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# Recoil Causes The Particle's Energy To Change, Scattered Photon Cannot Perform A "Strong" Measurement Of The Particles Position, Recoil-Heating Rate, Which Depends On The Photon-Scattering Rate And The Effect Of Each Scattering Event, Molecular Collisions Constitute The Dominant Noise Source, Photon-Recoil Heating Rate For Various Experimental Parameters And Quantum Predictions, Is E=Mc^2 An Exclusively Relativistic Result, Formula Should Be E= (E0)/(1-V^2/C^2)^(1/2) Derived By Laue In 1911 And Klein In 1918, Stochastic Gravitational-Wave Background From Binary Black Holes Analytical Computation Of The Mercury Perihelion Precession Via The Relativistic Gravitational Law And Comparison With General Relativity A Bohr-Type Model With Gravity As The Attractive Force

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**Abstract:** System disseminated and delineated upon is as follows: Recoil Causes The Particle's Energy To Change, Scattered Photon Cannot Perform A "Strong" Measurement Of The Particles Position, Recoil-Heating Rate, Which Depends On The Photon-Scattering Rate And The Effect Of Each Scattering Event, Molecular Collisions Constitute The Dominant Noise Source, Photon-Recoil Heating Rate For Various Experimental Parameters And Quantum Predictions, Is E=Mc^2 An Exclusively Relativistic Result, Formula Should Be E= (E0)/(1-V^2/C^2)^(1/2) Derived By Laue In 1911 And Klein In 1918, Stochastic Gravitational-Wave Background From Binary Black Holes, Analytical Computation Of The Mercury Perihelion Precession Via The Relativistic Gravitational Law And Comparison With General Relativity, A Bohr-Type Model With Gravity As The Attractive Force

**Key words**: Recoil Causes The Particle's Energy To Change, Scattered Photon, A "Strong" Measurement Of The Particles Position, Recoil-Heating Rate, Photon-Scattering Rate, Effect Of Each Scattering Event, Molecular Collisions Constitute The Dominant Noise Source, Photon-Recoil Heating Rate , Is E=Mc^2 An Exclusively Relativistic Result, Formula Should Be E= (E0)/(1-V^2/C^2)^(1/2) Derived By Laue In 1911 And Klein In 1918,

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**Introduction:**

System disseminated and delineated upon is as follows: Recoil Causes The Particle's Energy To Change, Scattered Photon Cannot Perform A "Strong" Measurement Of The Particles Position, Recoil-Heating Rate, Which Depends On The Photon-Scattering Rate And The Effect Of Each Scattering Event, Molecular Collisions Constitute The Dominant Noise Source, Photon-Recoil Heating Rate For Various Experimental Parameters And Quantum Predictions, Is E=Mc^2 An Exclusively Relativistic Result, Formula Should Be E= (E0)/(1-V^2/C^2)^(1/2) Derived By Laue In 1911 And Klein In 1918, Stochastic Gravitational-Wave Background From Binary Black Holes, Analytical Computation Of The Mercury Perihelion Precession Via The Relativistic Gravitational Law And Comparison With General Relativity, A Bohr-Type Model With Gravity As The Attractive Force

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## SECTION ONE

## Recoil Causes The Particle's Energy To Change

**INTRODUCTION—VARIABLES USED**

**Viewpoint: Measuring Quantum Kicks from a Beam of Light Amir Safavi-Naeini, 348 Via Pueblo Mall, Stanford University, Stanford, CA 94305-4088, USA June 13, 2016• Physics 9, 64 Force sensors levitated by light have reached the quantum regime, in which their sensitivity is limited by the momentum kicks of individual photons. Figure caption expand figure APS/Alan Stonebraker Figures and some sentences are deleted due to spatial restrictions Please excuse me**

(1) This recoil causes (eb) the particle's energy to (e) change by an amount of order of ℏ2k2⁄2mℏ2k2⁄2m, where (e) kk is the photon's wave number, mm is the particle's mass, and ℏℏ is the reduced Planck's constant.
(2) But in the ETH team's experiment, this recoil energy is nearly ten billion times smaller (eb) than the amount of energy needed to excite one quantum of the nanosphere's motion within (eb) the optical trap.
(3) An alternative way to understand the miniscule effect of (e) a single photon is to consider how much information about (e&eb) the trapped particle can be extracted from (e) the scattered photon.
(4) For the experimental parameters chosen by Novotny's group, the nanosphere's zero-point fluctuation length is (=) approximately 10 pm.
(5) This is the length scale on which a photon must measure (eb) the particle's position in order to discern (e&eb) quantum fluctuations. However, the photons from the trapping laser have a wavelength of 1 μμm, which is one hundred thousand times larger than (e) the fluctuation length.
(6) As such, a scattered photon cannot (e) perform a "strong" measurement of (e) the particle's position.
(7) In other words, the quantum back action from (e) the scattering of a single photon can have (e) only a minor effect on (e&eb) the dynamics of the particle.
(8) However, the cumulative momentum kicks from (e) many scattering events can change (e&eb) the

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| state of the nanosphere. | |
| **NOTATION** | |
| **Module One**<br><br>This recoil causes (eb) the particle's energy to (e) change by an amount of order of $\hbar$2k2⁄2m$\hbar$2k2⁄2m, where (e) kk is the photon's wave number, mm is the particle's mass, and $\hbar\hbar$ is the reduced Planck's constant | |
| $G_{13}$ : Category one of recoil ; particle's energy to (e) change by an amount of order of $\hbar$2k2⁄2m$\hbar$2k2⁄2m, where (e) kk is the photon's wave number, mm is the particle's mass, and $\hbar\hbar$ is the reduced Planck's constant<br><br>$G_{14}$ : Category two of SAS<br><br>$G_{15}$ : Category three of SAS | |
| $T_{13}$ : Category one of particle's energy to (e) change by an amount of order of $\hbar$2k2⁄2m$\hbar$2k2⁄2m, where (e) kk is the photon's wave number, mm is the particle's mass, and $\hbar\hbar$ is the reduced Planck's constant ;recoil<br><br>$T_{14}$ : Category two of SAS<br><br>$T_{15}$ : Category three of SAS | |
| **Module Two**<br><br>This recoil causes the particle's energy to (e) change by an amount of order of $\hbar$2k2⁄2m$\hbar$2k2⁄2m, where (e) kk is the photon's wave number, mm is the particle's mass, and $\hbar\hbar$ is the reduced Planck's constant | |
| $G_{16}$ : Category one of change by an amount of order of $\hbar$2k2⁄2m$\hbar$2k2⁄2m, where (e) kk is the photon's wave number, mm is the particle's mass, and $\hbar\hbar$ is the reduced Planck's constant<br><br>$G_{17}$ : Category two of SAS<br><br>$G_{18}$ : Category three of SAS | |
| $T_{16}$ : Category one of recoil causes the particle's energy<br><br>$T_{17}$ : Category two of SAS<br><br>$T_{18}$ : Category three of SAS | |
| **Module three**<br><br>This recoil causes the particle's energy to change by an amount of order of $\hbar$2k2⁄2m$\hbar$2k2⁄2m, where (e) kk is the photon's wave number, mm is the particle's mass, and $\hbar\hbar$ is the reduced Planck's constant | |
| $G_{20}$ : Category one of kk is the photon's wave number, mm is the particle's mass, and $\hbar\hbar$ is the reduced Planck's constant<br><br>$G_{21}$ : Category two of SAS | |

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| $G_{22}$ : Category three of SAS | |
| $T_{20}$ : Category one of recoil causes the particle's energy to change by an amount of order of ħ2k2/2mħ2k2/2m<br><br>$T_{21}$ : Category two of SAS<br><br>$T_{22}$ : Category three of SAS | |
| **Module four**<br><br>But in the ETH team's experiment, this recoil energy is nearly ten billion times smaller (eb) than the amount of energy needed to excite one quantum of the nanosphere's motion within (eb) the optical trap | |
| $G_{24}$ : Category one of ETH team's experiment, this recoil energy is nearly ten billion times<br><br>$G_{25}$ : Category two of SAS<br><br>$G_{26}$ : Category three of SAS | |
| $T_{24}$ : Category one of amount of energy needed to excite one quantum of the nanosphere's motion within (eb) the optical trap<br><br>$T_{25}$ : Category two of SAS<br><br>$T_{26}$ : Category three of SAS | |
| **Module five**<br><br>But, this recoil energy in the ETH team's experiment is nearly ten billion times smaller (eb) than the amount of energy needed to excite one quantum of the nanosphere's motion within (eb) the optical trap | |
| $G_{28}$ : Category one of this recoil energy in the ETH team's experiment is nearly ten billion times smaller<br><br>$G_{29}$ : Category two of SAS<br><br>$G_{30}$ : Category three of SAS | |
| $T_{28}$ : Category one of amount of energy needed to excite one quantum of the nanosphere's motion within (eb) the optical trap<br><br>$T_{29}$ : Category two of SAS<br><br>$T_{30}$ : Category three of SAS | |
| **Module six**<br><br>But, this recoil energy in the ETH team's experiment is nearly ten billion times smaller than the amount of energy needed to excite one quantum of the nanosphere's motion within (eb) the optical trap | |
| $G_{32}$ : Category one of recoil energy in the ETH team's experiment is nearly ten billion times smaller than | |

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| the amount of energy needed to excite one quantum of the nanosphere's motion<br><br>$G_{33}$ : Category two of SAS<br><br>$G_{34}$ : Category three of SAS | |
| $T_{32}$ : Category one of optical trap (S)<br><br>$T_{33}$ : Category two of SAS<br><br>$T_{34}$ : Category three of SAS | |
| **Module seven**<br><br>An alternative way to understand the miniscule effect of (e) a single photon is to consider how much information about (e&eb) the trapped particle can be extracted from (e) the scattered photon | |
| $G_{36}$ : Category one of single photon is to consider how much information about (e&eb) the trapped particle can be extracted from (e) the scattered photon<br><br>$G_{37}$ : Category two of SAS<br><br>$G_{38}$ : Category three of SAS | |
| $T_{36}$ : Category one of alternative way to understand the miniscule effect<br><br>$T_{37}$ : Category two of SAS<br><br>$T_{38}$ : Category three of SAS | |
| **Module eight**<br><br>An alternative way to understand the miniscule effect of a single photon is to consider how much information about (e&eb) the trapped particle can be extracted from (e) the scattered photon | |
| $G_{40}$ : Category one of **alternative way to understand the miniscule effect of a single photon is to consider how much information;** the trapped particle can be extracted from (e) the scattered photon<br><br>$G_{41}$ : Category two of SAS<br><br>$G_{42}$ : Category three of SAS | |
| $T_{40}$ : Category one of the trapped particle can be extracted from (e) the scattered photon **;alternative way to understand the miniscule effect of a single photon is to consider how much information**<br><br>$T_{41}$ : Category two of SAS<br><br>$T_{42}$ : Category three of SAS | |
| **Module Nine**<br><br>An alternative way to understand the miniscule effect of a single photon is to consider how much | |

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| information about the trapped particle can be extracted from (e) the scattered photon | |
| $G_{44}$ : Category one of scattered photon<br><br>$G_{45}$ : Category two of SAS<br><br>$G_{46}$ : Category three of SAS | |
| $T_{44}$ : Category one of alternative way to understand the miniscule effect of a single photon is to consider how much information about the trapped particle can be extracted<br><br>$T_{45}$ : Category two of SAS<br><br>$T_{46}$ : Category three of SAS | |
| **The Coefficients:** | |
| $(a_{13})^{(1)}, (a_{14})^{(1)}, (a_{15})^{(1)}, (b_{13})^{(1)}, (b_{14})^{(1)}, (b_{15})^{(1)}\ (a_{16})^{(2)}, (a_{17})^{(2)}, (a_{18})^{(2)}\ (b_{16})^{(2)}, (b_{17})^{(2)}, (b_{18})^{(2)}$:<br>$(a_{20})^{(3)}, (a_{21})^{(3)}, (a_{22})^{(3)}$ , $(b_{20})^{(3)}, (b_{21})^{(3)}, (b_{22})^{(3)}$<br>$(a_{24})^{(4)}, (a_{25})^{(4)}, (a_{26})^{(4)}, (b_{24})^{(4)}, (b_{25})^{(4)}, (b_{26})^{(4)}, (b_{28})^{(5)}, (b_{29})^{(5)}, (b_{30})^{(5)}, (a_{28})^{(5)}, (a_{29})^{(5)}, (a_{30})^{(5)}$,<br>$(a_{32})^{(6)}, (a_{33})^{(6)}, (a_{34})^{(6)}, (b_{32})^{(6)}, (b_{33})^{(6)}, (b_{34})^{(6)}$<br>$(a_{36})^{(7)}, (a_{37})^{(7)}, (a_{38})^{(7)}, (b_{36})^{(7)}, (b_{37})^{(7)}, (b_{38})^{(7)}$<br>$(a_{40})^{(8)}, (a_{41})^{(8)}, (a_{42})^{(8)}, (b_{40})^{(8)}, (b_{41})^{(8)}, (b_{42})^{(8)}$<br>$(a_{44})^{(9)}, (a_{45})^{(9)}, (a_{46})^{(9)}, (b_{44})^{(9)}, (b_{45})^{(9)}, (b_{46})^{(9)}$<br><br>are Accentuation coefficients<br>$(a_{13}^{'})^{(1)}, (a_{14}^{'})^{(1)}, (a_{15}^{'})^{(1)}, (b_{13}^{'})^{(1)}, (b_{14}^{'})^{(1)}, (b_{15}^{'})^{(1)}, (a_{16}^{'})^{(2)}, (a_{17}^{'})^{(2)}, (a_{18}^{'})^{(2)}$, $(b_{16}^{'})^{(2)}, (b_{17}^{'})^{(2)}, (b_{18}^{'})^{(2)}$<br>$, (a_{20}^{'})^{(3)}, (a_{21}^{'})^{(3)}, (a_{22}^{'})^{(3)}, (b_{20}^{'})^{(3)}, (b_{21}^{'})^{(3)}, (b_{22}^{'})^{(3)}$<br>$(a_{24}^{'})^{(4)}, (a_{25}^{'})^{(4)}, (a_{26}^{'})^{(4)}, (b_{24}^{'})^{(4)}, (b_{25}^{'})^{(4)}, (b_{26}^{'})^{(4)}, (b_{28}^{'})^{(5)}, (b_{29}^{'})^{(5)}, (b_{30}^{'})^{(5)}\ (a_{28}^{'})^{(5)}, (a_{29}^{'})^{(5)}, (a_{30}^{'})^{(5)}$ ,<br>$(a_{32}^{'})^{(6)}, (a_{33}^{'})^{(6)}, (a_{34}^{'})^{(6)}, (b_{32}^{'})^{(6)}, (b_{33}^{'})^{(6)}, (b_{34}^{'})^{(6)}$<br>$(a_{36}^{'})^{(7)}, (a_{37}^{'})^{(7)}, (a_{38}^{'})^{(7)}, (b_{36}^{'})^{(7)}, (b_{37}^{'})^{(7)}, (b_{38}^{'})^{(7)}$,<br>$(a_{40}^{'})^{(8)}, (a_{41}^{'})^{(8)}, (a_{42}^{'})^{(8)}, (b_{40}^{'})^{(8)}, (b_{41}^{'})^{(8)}, (b_{42}^{'})^{(8)}$,<br>$(a_{44}^{'})^{(9)}, (a_{45}^{'})^{(9)}, (a_{46}^{'})^{(9)}, (b_{44}^{'})^{(9)}, (b_{45}^{'})^{(9)}, (b_{46}^{'})^{(9)}$,<br>are Dissipation coefficients | |
| **Module Numbered One**<br><br>**The differential system of this model is now (Module Numbered one)** | |
| $\frac{dG_{13}}{dt} = (a_{13})^{(1)}G_{14} - \left[(a_{13}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14}, t)\right]G_{13}$ | 1 |
| $\frac{dG_{14}}{dt} = (a_{14})^{(1)}G_{13} - \left[(a_{14}^{'})^{(1)} + (a_{14}^{''})^{(1)}(T_{14}, t)\right]G_{14}$ | 2 |
| $\frac{dG_{15}}{dt} = (a_{15})^{(1)}G_{14} - \left[(a_{15}^{'})^{(1)} + (a_{15}^{''})^{(1)}(T_{14}, t)\right]G_{15}$ | 3 |
| $\frac{dT_{13}}{dt} = (b_{13})^{(1)}T_{14} - \left[(b_{13}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G, t)\right]T_{13}$ | 4 |
| $\frac{dT_{14}}{dt} = (b_{14})^{(1)}T_{13} - \left[(b_{14}^{'})^{(1)} - (b_{14}^{''})^{(1)}(G, t)\right]T_{14}$ | 5 |
| $\frac{dT_{15}}{dt} = (b_{15})^{(1)}T_{14} - \left[(b_{15}^{'})^{(1)} - (b_{15}^{''})^{(1)}(G, t)\right]T_{15}$ | 6 |
| $+(a_{13}^{''})^{(1)}(T_{14}, t) =$ First augmentation factor | |

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| $-(b_{13}^{''})^{(1)}(G,t) =$ First detritions factor | |
| **Module Numbered Two**<br><br>**The differential system of this model is now ( Module numbered two)** | |
| $\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \left[(a_{16}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17},t)\right]G_{16}$ | 7 |
| $\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \left[(a_{17}^{'})^{(2)} + (a_{17}^{''})^{(2)}(T_{17},t)\right]G_{17}$ | 8 |
| $\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \left[(a_{18}^{'})^{(2)} + (a_{18}^{''})^{(2)}(T_{17},t)\right]G_{18}$ | 9 |
| $\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \left[(b_{16}^{'})^{(2)} - (b_{16}^{''})^{(2)}((G_{19}),t)\right]T_{16}$ | 10 |
| $\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \left[(b_{17}^{'})^{(2)} - (b_{17}^{''})^{(2)}((G_{19}),t)\right]T_{17}$ | 11 |
| $\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \left[(b_{18}^{'})^{(2)} - (b_{18}^{''})^{(2)}((G_{19}),t)\right]T_{18}$ | 12 |
| $+(a_{16}^{''})^{(2)}(T_{17},t) =$ First augmentation factor | |
| $-(b_{16}^{''})^{(2)}((G_{19}),t) =$ First detritions factor | |
| **Module Numbered Three**<br><br>**The differential system of this model is now (Module numbered three)** | |
| $\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \left[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21},t)\right]G_{20}$ | 13 |
| $\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \left[(a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)}(T_{21},t)\right]G_{21}$ | 14 |
| $\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \left[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21},t)\right]G_{22}$ | 15 |
| $\frac{dT_{20}}{dt} = (b_{20})^{(3)}T_{21} - \left[(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23},t)\right]T_{20}$ | 16 |
| $\frac{dT_{21}}{dt} = (b_{21})^{(3)}T_{20} - \left[(b_{21}^{'})^{(3)} - (b_{21}^{''})^{(3)}(G_{23},t)\right]T_{21}$ | 17 |
| $\frac{dT_{22}}{dt} = (b_{22})^{(3)}T_{21} - \left[(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23},t)\right]T_{22}$ | 18 |
| $+(a_{20}^{''})^{(3)}(T_{21},t) =$ First augmentation factor | |
| $-(b_{20}^{''})^{(3)}(G_{23},t) =$ First detritions factor | |
| **Module Numbered Four**<br><br>**The differential system of this model is now (Module numbered Four)** | |
| $\frac{dG_{24}}{dt} = (a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25},t)\right]G_{24}$ | 19 |
| $\frac{dG_{25}}{dt} = (a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25},t)\right]G_{25}$ | 20 |
| $\frac{dG_{26}}{dt} = (a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25},t)\right]G_{26}$ | 21 |
| $\frac{dT_{24}}{dt} = (b_{24})^{(4)}T_{25} - \left[(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}((G_{27}),t)\right]T_{24}$ | 22 |
| $\frac{dT_{25}}{dt} = (b_{25})^{(4)}T_{24} - \left[(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}((G_{27}),t)\right]T_{25}$ | 23 |
| $\frac{dT_{26}}{dt} = (b_{26})^{(4)}T_{25} - \left[(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}((G_{27}),t)\right]T_{26}$ | 24 |
| $+(a_{24}^{''})^{(4)}(T_{25},t) =$ **First augmentation factor** | |
| $-(b_{24}^{''})^{(4)}((G_{27}),t) =$ **First detritions factor** | |
| **Module Numbered Five**:<br><br>**The differential system of this model is now (Module number five)** | |

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| $\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}, t)\right]G_{28}$ | 25 |
| $\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29}, t)\right]G_{29}$ | 26 |
| $\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \left[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29}, t)\right]G_{30}$ | 27 |
| $\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \left[(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}((G_{31}), t)\right]T_{28}$ | 28 |
| $\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \left[(b_{29}^{'})^{(5)} - (b_{29}^{''})^{(5)}((G_{31}), t)\right]T_{29}$ | 29 |
| $\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \left[(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}((G_{31}), t)\right]T_{30}$ | 30 |
| $+(a_{28}^{''})^{(5)}(T_{29}, t) =$ **First augmentation factor** | |
| $-(b_{28}^{''})^{(5)}((G_{31}), t) =$ **First detritions factor** | |
| **Module Numbered Six**<br><br>**The differential system of this model is now (Module numbered Six)** | |
| $\frac{dG_{32}}{dt} = (a_{32})^{(6)}G_{33} - \left[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33}, t)\right]G_{32}$ | 31 |
| $\frac{dG_{33}}{dt} = (a_{33})^{(6)}G_{32} - \left[(a_{33}^{'})^{(6)} + (a_{33}^{''})^{(6)}(T_{33}, t)\right]G_{33}$ | 32 |
| $\frac{dG_{34}}{dt} = (a_{34})^{(6)}G_{33} - \left[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33}, t)\right]G_{34}$ | 33 |
| $\frac{dT_{32}}{dt} = (b_{32})^{(6)}T_{33} - \left[(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}((G_{35}), t)\right]T_{32}$ | 34 |
| $\frac{dT_{33}}{dt} = (b_{33})^{(6)}T_{32} - \left[(b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}((G_{35}), t)\right]T_{33}$ | 35 |
| $\frac{dT_{34}}{dt} = (b_{34})^{(6)}T_{33} - \left[(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}((G_{35}), t)\right]T_{34}$ | 36 |
| $+(a_{32}^{''})^{(6)}(T_{33}, t) =$ **First augmentation factor** | |
| **Module Numbered Seven:**<br><br>**The differential system of this model is now (Seventh Module)** | |
| $\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \left[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37}, t)\right]G_{36}$ | 37 |
| $\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \left[(a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37}, t)\right]G_{37}$ | *38* |
| $\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \left[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37}, t)\right]G_{38}$ | *39* |
| $\frac{dT_{36}}{dt} = (b_{36})^{(7)}T_{37} - \left[(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}((G_{39}), t)\right]T_{36}$ | *40* |
| $\frac{dT_{37}}{dt} = (b_{37})^{(7)}T_{36} - \left[(b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}((G_{39}), t)\right]T_{37}$ | *41* |
| $\frac{dT_{38}}{dt} = (b_{38})^{(7)}T_{37} - \left[(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}((G_{39}), t)\right]T_{38}$ | *42* |
| $+(a_{36}^{''})^{(7)}(T_{37}, t) =$ **First augmentation factor** | |
| **Module Numbered Eight**<br><br>The differential system of this model is now | |
| $\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \left[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41}, t)\right]G_{40}$ | *43* |
| $\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \left[(a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41}, t)\right]G_{41}$ | *44* |
| $\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \left[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41}, t)\right]G_{42}$ | *45* |
| $\frac{dT_{40}}{dt} = (b_{40})^{(8)}T_{41} - \left[(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}((G_{43}), t)\right]T_{40}$ | *46* |
| $\frac{dT_{41}}{dt} = (b_{41})^{(8)}T_{40} - \left[(b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}((G_{43}), t)\right]T_{41}$ | *47* |

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| $\frac{dT_{42}}{dt} = (b_{42})^{(8)}T_{41} - \left[(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}((G_{43}),t)\right]T_{42}$ | 48 |
| **Module Numbered Nine**<br><br>The differential system of this model is now | |
| $\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \left[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45},t)\right]G_{44}$ | 49 |
| $\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \left[(a_{45}^{'})^{(9)} + (a_{45}^{''})^{(9)}(T_{45},t)\right]G_{45}$ | 50 |
| $\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \left[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45},t)\right]G_{46}$ | 51 |
| $\frac{dT_{44}}{dt} = (b_{44})^{(9)}T_{45} - \left[(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}((G_{47}),t)\right]T_{44}$ | 52 |
| $\frac{dT_{45}}{dt} = (b_{45})^{(9)}T_{44} - \left[(b_{45}^{'})^{(9)} - (b_{45}^{''})^{(9)}((G_{47}),t)\right]T_{45}$ | 53 |
| $\frac{dT_{46}}{dt} = (b_{46})^{(9)}T_{45} - \left[(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}((G_{47}),t)\right]T_{46}$ | 54 |
| $+(a_{44}^{''})^{(9)}(T_{45},t) =$ **First augmentation factor** | |
| $-(b_{44}^{''})^{(9)}((G_{47}),t) =$ **First detrition factor** | |
| $\frac{dG_{13}}{dt} = (a_{13})^{(1)}G_{14} - \left[\begin{array}{ccc} (a_{13}^{'})^{(1)}\boxed{+(a_{13}^{''})^{(1)}(T_{14},t)} & \boxed{+(a_{16}^{''})^{(2,2,)}(T_{17},t)} & \boxed{+(a_{20}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{24}^{''})^{(4,4,4,4,)}(T_{25},t)} & \boxed{+(a_{28}^{''})^{(5,5,5,5,)}(T_{29},t)} & \boxed{+(a_{32}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{36}^{''})^{(7,7)}(T_{37},t)} & \boxed{+(a_{40}^{''})^{(8,8)}(T_{41},t)} & \boxed{+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{array}\right]G_{13}$ | 55 |
| $\frac{dG_{14}}{dt} = (a_{14})^{(1)}G_{13} - \left[\begin{array}{ccc} (a_{14}^{'})^{(1)}\boxed{+(a_{14}^{''})^{(1)}(T_{14},t)} & \boxed{+(a_{17}^{''})^{(2,2,)}(T_{17},t)} & \boxed{+(a_{21}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{25}^{''})^{(4,4,4,4,)}(T_{25},t)} & \boxed{+(a_{29}^{''})^{(5,5,5,5,)}(T_{29},t)} & \boxed{+(a_{33}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{37}^{''})^{(7,7)}(T_{37},t)} & \boxed{+(a_{41}^{''})^{(8,8)}(T_{41},t)} & \boxed{+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{array}\right]G_{14}$ | 56 |
| $\frac{dG_{15}}{dt} = (a_{15})^{(1)}G_{14} - \left[\begin{array}{ccc} (a_{15}^{'})^{(1)}\boxed{+(a_{15}^{''})^{(1)}(T_{14},t)} & \boxed{+(a_{18}^{''})^{(2,2,)}(T_{17},t)} & \boxed{+(a_{22}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{26}^{''})^{(4,4,4,4,)}(T_{25},t)} & \boxed{+(a_{30}^{''})^{(5,5,5,5,)}(T_{29},t)} & \boxed{+(a_{34}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{38}^{''})^{(7,7)}(T_{37},t)} & \boxed{+(a_{42}^{''})^{(8,8)}(T_{41},t)} & \boxed{+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{array}\right]G_{15}$ | 57 |
| Where $\boxed{(a_{13}^{''})^{(1)}(T_{14},t)}$, $\boxed{(a_{14}^{''})^{(1)}(T_{14},t)}$, $\boxed{(a_{15}^{''})^{(1)}(T_{14},t)}$ are first augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{16}^{''})^{(2,2,)}(T_{17},t)}$, $\boxed{+(a_{17}^{''})^{(2,2,)}(T_{17},t)}$, $\boxed{+(a_{18}^{''})^{(2,2,)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{20}^{''})^{(3,3,)}(T_{21},t)}$, $\boxed{+(a_{21}^{''})^{(3,3,)}(T_{21},t)}$, $\boxed{+(a_{22}^{''})^{(3,3,)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{24}^{''})^{(4,4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{25}^{''})^{(4,4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{26}^{''})^{(4,4,4,4,)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{28}^{''})^{(5,5,5,5,)}(T_{29},t)}$, $\boxed{+(a_{29}^{''})^{(5,5,5,5,)}(T_{29},t)}$, $\boxed{+(a_{30}^{''})^{(5,5,5,5,)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{32}^{''})^{(6,6,6,6,)}(T_{33},t)}$, $\boxed{+(a_{33}^{''})^{(6,6,6,6,)}(T_{33},t)}$, $\boxed{+(a_{34}^{''})^{(6,6,6,6,)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{38}^{''})^{(7,7)}(T_{37},t)}$ $\boxed{+(a_{37}^{''})^{(7,7)}(T_{37},t)}$ $\boxed{+(a_{36}^{''})^{(7,7)}(T_{37},t)}$ are seventh augmentation coefficient for 1,2,3<br>$\boxed{+(a_{40}^{''})^{(8,8)}(T_{41},t)}$ $\boxed{+(a_{41}^{''})^{(8,8)}(T_{41},t)}$ $\boxed{+(a_{42}^{''})^{(8,8)}(T_{41},t)}$ are eight augmentation coefficient for 1,2,3 | |

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| $+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficient for 1,2,3 | |
| $\frac{dT_{13}}{dt} = (b_{13})^{(1)}T_{14} - \begin{bmatrix} (b_{13}^{'})^{(1)}\left[-(b_{13}^{''})^{(1)}(G,t)\right] & \left[-(b_{16}^{''})^{(2,2,)}(G_{19},t)\right] & \left[-(b_{20}^{''})^{(3,3,)}(G_{23},t)\right] \\ \left[-(b_{24}^{''})^{(4,4,4,4,)}(G_{27},t)\right] & \left[-(b_{28}^{''})^{(5,5,5,5,)}(G_{31},t)\right] & \left[-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)\right] \\ \left[-(b_{36}^{''})^{(7,7,)}(G_{39},t)\right] & \left[-(b_{40}^{''})^{(8,8)}(G_{43},t)\right] & \left[-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{13}$ | 58 |
| $\frac{dT_{14}}{dt} = (b_{14})^{(1)}T_{13} - \begin{bmatrix} (b_{14}^{'})^{(1)}\left[-(b_{14}^{''})^{(1)}(G,t)\right] & \left[-(b_{17}^{''})^{(2,2,)}(G_{19},t)\right] & \left[-(b_{21}^{''})^{(3,3,)}(G_{23},t)\right] \\ \left[-(b_{25}^{''})^{(4,4,4,4,)}(G_{27},t)\right] & \left[-(b_{29}^{''})^{(5,5,5,5,)}(G_{31},t)\right] & \left[-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)\right] \\ \left[-(b_{37}^{''})^{(7,7,)}(G_{39},t)\right] & \left[-(b_{41}^{''})^{(8,8)}(G_{43},t)\right] & \left[-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{14}$ | 59 |
| $\frac{dT_{15}}{dt} = (b_{15})^{(1)}T_{14} - \begin{bmatrix} (b_{15}^{'})^{(1)}\left[-(b_{15}^{''})^{(1)}(G,t)\right] & \left[-(b_{18}^{''})^{(2,2,)}(G_{19},t)\right] & \left[-(b_{22}^{''})^{(3,3,)}(G_{23},t)\right] \\ \left[-(b_{26}^{''})^{(4,4,4,4,)}(G_{27},t)\right] & \left[-(b_{30}^{''})^{(5,5,5,5,)}(G_{31},t)\right] & \left[-(b_{34}^{''})^{(6,6,6,6,)}(G_{35},t)\right] \\ \left[-(b_{38}^{''})^{(7,7,)}(G_{39},t)\right] & \left[-(b_{42}^{''})^{(8,8)}(G_{43},t)\right] & \left[-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{15}$ | *60* |
| Where $\left[-(b_{13}^{''})^{(1)}(G,t)\right]$, $\left[-(b_{14}^{''})^{(1)}(G,t)\right]$, $\left[-(b_{15}^{''})^{(1)}(G,t)\right]$ are first detrition coefficients for category 1, 2 and 3<br>$\left[-(b_{16}^{''})^{(2,2,)}(G_{19},t)\right]$, $\left[-(b_{17}^{''})^{(2,2,)}(G_{19},t)\right]$, $\left[-(b_{18}^{''})^{(2,2,)}(G_{19},t)\right]$ are second detrition coefficients for category 1, 2 and 3<br>$\left[-(b_{20}^{''})^{(3,3,)}(G_{23},t)\right]$, $\left[-(b_{21}^{''})^{(3,3,)}(G_{23},t)\right]$, $\left[-(b_{22}^{''})^{(3,3,)}(G_{23},t)\right]$ are third detrition coefficients for category 1, 2 and 3<br>$\left[-(b_{24}^{''})^{(4,4,4,4,)}(G_{27},t)\right]$, $\left[-(b_{25}^{''})^{(4,4,4,4,)}(G_{27},t)\right]$, $\left[-(b_{26}^{''})^{(4,4,4,4,)}(G_{27},t)\right]$ are fourth detrition coefficients for category 1, 2 and 3<br>$\left[-(b_{28}^{''})^{(5,5,5,5,)}(G_{31},t)\right]$, $\left[-(b_{29}^{''})^{(5,5,5,5,)}(G_{31},t)\right]$, $\left[-(b_{30}^{''})^{(5,5,5,5,)}(G_{31},t)\right]$ are fifth detrition coefficients for category 1, 2 and 3<br>$\left[-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)\right]$, $\left[-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)\right]$, $\left[-(b_{34}^{''})^{(6,6,6,6,)}(G_{35},t)\right]$ are sixth detrition coefficients for category 1, 2 and 3<br>$\left[-(b_{37}^{''})^{(7,7,)}(G_{39},t)\right]$, $\left[-(b_{36}^{''})^{(7,7,)}(G_{39},t)\right]$, $\left[-(b_{38}^{''})^{(7,7,)}(G_{39},t)\right]$ are seventh detrition coefficients for category 1, 2 and 3<br>$\left[-(b_{40}^{''})^{(8,8)}(G_{43},t)\right]$ $\left[-(b_{41}^{''})^{(8,8)}(G_{43},t)\right]$ $\left[-(b_{42}^{''})^{(8,8)}(G_{43},t)\right]$ are eight detrition coefficients for category 1, 2 and 3<br>$\left[-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right]$, $\left[-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right]$, $\left[-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right]$ are ninth detrition coefficients for category 1, 2 and 3 | |
| | |
| $\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \begin{bmatrix} (a_{16}^{'})^{(2)}\left[+(a_{16}^{''})^{(2)}(T_{17},t)\right] & \left[+(a_{13}^{''})^{(1,1,)}(T_{14},t)\right] & \left[+(a_{20}^{''})^{(3,3,3)}(T_{21},t)\right] \\ \left[+(a_{24}^{''})^{(4,4,4,4,4)}(T_{25},t)\right] & \left[+(a_{28}^{''})^{(5,5,5,5,5)}(T_{29},t)\right] & \left[+(a_{32}^{''})^{(6,6,6,6,6)}(T_{33},t)\right] \\ \left[+(a_{36}^{''})^{(7,7,7)}(T_{37},t)\right] & \left[+(a_{40}^{''})^{(8,8,8)}(T_{41},t)\right] & \left[+(a_{44}^{''})^{(9,9)}(T_{45},t)\right] \end{bmatrix} G_{16}$ | *61* |

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| $\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \left[\begin{array}{ccc} (a_{17}^{'})^{(2)}\boxed{+(a_{17}^{''})^{(2)}(T_{17},t)} & \boxed{+(a_{14}^{''})^{(1,1,)}(T_{14},t)} & \boxed{+(a_{21}^{''})^{(3,3,3)}(T_{21},t)} \\ \boxed{+(a_{25}^{''})^{(4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{29}^{''})^{(5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{33}^{''})^{(6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{37}^{''})^{(7,7,7)}(T_{37},t)} & \boxed{+(a_{41}^{''})^{(8,8,8)}(T_{41},t)} & \boxed{+(a_{45}^{''})^{(9,9)}(T_{45},t)} \end{array}\right] G_{17}$ | 62 |
| $\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \left[\begin{array}{ccc} (a_{18}^{'})^{(2)}\boxed{+(a_{18}^{''})^{(2)}(T_{17},t)} & \boxed{+(a_{15}^{''})^{(1,1,)}(T_{14},t)} & \boxed{+(a_{22}^{''})^{(3,3,3)}(T_{21},t)} \\ \boxed{+(a_{26}^{''})^{(4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{30}^{''})^{(5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{34}^{''})^{(6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{38}^{''})^{(7,7,7)}(T_{37},t)} & \boxed{+(a_{42}^{''})^{(8,8,8)}(T_{41},t)} & \boxed{+(a_{46}^{''})^{(9,9)}(T_{45},t)} \end{array}\right] G_{18}$ | 63 |
| Where $\boxed{+(a_{16}^{''})^{(2)}(T_{17},t)}$, $\boxed{+(a_{17}^{''})^{(2)}(T_{17},t)}$, $\boxed{+(a_{18}^{''})^{(2)}(T_{17},t)}$ are first augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{13}^{''})^{(1,1,)}(T_{14},t)}$, $\boxed{+(a_{14}^{''})^{(1,1,)}(T_{14},t)}$, $\boxed{+(a_{15}^{''})^{(1,1,)}(T_{14},t)}$ are second augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{20}^{''})^{(3,3,3)}(T_{21},t)}$, $\boxed{+(a_{21}^{''})^{(3,3,3)}(T_{21},t)}$, $\boxed{+(a_{22}^{''})^{(3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{24}^{''})^{(4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{25}^{''})^{(4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{26}^{''})^{(4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{28}^{''})^{(5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{29}^{''})^{(5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{30}^{''})^{(5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{32}^{''})^{(6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{33}^{''})^{(6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{34}^{''})^{(6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{36}^{''})^{(7,7,7)}(T_{37},t)}$, $\boxed{+(a_{37}^{''})^{(7,7,7)}(T_{37},t)}$, $\boxed{+(a_{38}^{''})^{(7,7,7)}(T_{37},t)}$ are seventh augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{40}^{''})^{(8,8,8)}(T_{41},t)}$, $\boxed{+(a_{41}^{''})^{(8,8,8)}(T_{41},t)}$, $\boxed{+(a_{42}^{''})^{(8,8,8)}(T_{41},t)}$ are eight augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{44}^{''})^{(9,9)}(T_{45},t)}$, $\boxed{+(a_{45}^{''})^{(9,9)}(T_{45},t)}$, $\boxed{+(a_{46}^{''})^{(9,9)}(T_{45},t)}$ are ninth augmentation coefficient for category 1, 2 and 3 | |
| $\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \left[\begin{array}{ccc} (b_{16}^{'})^{(2)}\boxed{-(b_{16}^{''})^{(2)}(G_{19},t)} & \boxed{-(b_{13}^{''})^{(1,1,)}(G,t)} & \boxed{-(b_{20}^{''})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b_{24}^{''})^{(4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{28}^{''})^{(5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{32}^{''})^{(6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{36}^{''})^{(7,7,7)}(G_{39},t)} & \boxed{-(b_{40}^{''})^{(8,8,8)}(G_{43},t)} & \boxed{-(b_{44}^{''})^{(9,9)}(G_{47},t)} \end{array}\right] T_{16}$ | 64 |
| $\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \left[\begin{array}{ccc} (b_{17}^{'})^{(2)}\boxed{-(b_{17}^{''})^{(2)}(G_{19},t)} & \boxed{-(b_{14}^{''})^{(1,1,)}(G,t)} & \boxed{-(b_{21}^{''})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b_{25}^{''})^{(4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{29}^{''})^{(5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{33}^{''})^{(6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{37}^{''})^{(7,7,7)}(G_{39},t)} & \boxed{-(b_{41}^{''})^{(8,8,8)}(G_{43},t)} & \boxed{-(b_{45}^{''})^{(9,9)}(G_{47},t)} \end{array}\right] T_{17}$ | 65 |
| $\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \left[\begin{array}{ccc} (b_{18}^{'})^{(2)}\boxed{-(b_{18}^{''})^{(2)}(G_{19},t)} & \boxed{-(b_{15}^{''})^{(1,1,)}(G,t)} & \boxed{-(b_{22}^{''})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b_{26}^{''})^{(4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{30}^{''})^{(5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{34}^{''})^{(6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{38}^{''})^{(7,7,7)}(G_{39},t)} & \boxed{-(b_{42}^{''})^{(8,8,8)}(G_{43},t)} & \boxed{-(b_{46}^{''})^{(9,9)}(G_{47},t)} \end{array}\right] T_{18}$ | 66 |
| where $\boxed{-(b_{16}^{''})^{(2)}(G_{19},t)}$, $\boxed{-(b_{17}^{''})^{(2)}(G_{19},t)}$, $\boxed{-(b_{18}^{''})^{(2)}(G_{19},t)}$ are first detrition coefficients for category 1, 2 and 3 | |

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| $-(b_{13}^{''})^{(1,1,)}(G,t)$, $-(b_{14}^{''})^{(1,1,)}(G,t)$, $-(b_{15}^{''})^{(1,1,)}(G,t)$ are second detrition coefficients for category 1,2 and 3<br>$-(b_{20}^{''})^{(3,3,3,)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,)}(G_{23},t)$ are third detrition coefficients for category 1,2 and 3<br>$-(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1,2 and 3<br>$-(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1,2 and 3<br>$-(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$ are sixth detrition coefficients for category 1,2 and 3<br>$-(b_{36}^{''})^{(7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1,2 and 3<br>$-(b_{40}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8)}(G_{43},t)$ are eight detrition coefficients for category 1,2 and 3<br>$-(b_{44}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1,2 and 3 | |
| $\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \left[ \begin{array}{c} (a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21},t) + (a_{16}^{''})^{(2,2,2)}(T_{17},t) + (a_{13}^{''})^{(1,1,1,)}(T_{14},t) \\ + (a_{24}^{''})^{(4,4,4,4,4,4)}(T_{25},t) + (a_{28}^{''})^{(5,5,5,5,5,5)}(T_{29},t) + (a_{32}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ + (a_{36}^{''})^{(7,7,7,7)}(T_{37},t) + (a_{40}^{''})^{(8,8,8,8)}(T_{41},t) + (a_{44}^{''})^{(9,9,9)}(T_{45},t) \end{array} \right] G_{20}$ | 67 |
| $\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \left[ \begin{array}{c} (a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)}(T_{21},t) + (a_{17}^{''})^{(2,2,2)}(T_{17},t) + (a_{14}^{''})^{(1,1,1,)}(T_{14},t) \\ + (a_{25}^{''})^{(4,4,4,4,4,4)}(T_{25},t) + (a_{29}^{''})^{(5,5,5,5,5,5)}(T_{29},t) + (a_{33}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ + (a_{37}^{''})^{(7,7,7,7)}(T_{37},t) + (a_{41}^{''})^{(8,8,8,8)}(T_{41},t) + (a_{45}^{''})^{(9,9,9)}(T_{45},t) \end{array} \right] G_{21}$ | 68 |
| $\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \left[ \begin{array}{c} (a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21},t) + (a_{18}^{''})^{(2,2,2)}(T_{17},t) + (a_{15}^{''})^{(1,1,1,)}(T_{14},t) \\ + (a_{26}^{''})^{(4,4,4,4,4,4)}(T_{25},t) + (a_{30}^{''})^{(5,5,5,5,5,5)}(T_{29},t) + (a_{34}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ + (a_{38}^{''})^{(7,7,7,7)}(T_{37},t) + (a_{42}^{''})^{(8,8,8,8)}(T_{41},t) + (a_{46}^{''})^{(9,9,9)}(T_{45},t) \end{array} \right] G_{22}$ | 69 |
| $+(a_{20}^{''})^{(3)}(T_{21},t)$, $+(a_{21}^{''})^{(3)}(T_{21},t)$, $+(a_{22}^{''})^{(3)}(T_{21},t)$ are first augmentation coefficients for category 1, 2 and 3<br>$+(a_{16}^{''})^{(2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2)}(T_{17},t)$ are second augmentation coefficients for category 1, 2 and 3<br>$+(a_{13}^{''})^{(1,1,1,)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,)}(T_{14},t)$ are third augmentation coefficients for category 1, 2 and 3<br>$+(a_{24}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficients for category 1, 2 and 3<br>$+(a_{28}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficients for category 1, 2 and 3<br>$+(a_{32}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation | |

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| coefficients for category 1, 2 and 3 $+(a_{36}^{''})^{(7,7,7,7)}(T_{37},t)$, $+(a_{37}^{''})^{(7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7)}(T_{37},t)$ are seventh augmentation coefficients for category 1, 2 and 3 $+(a_{40}^{''})^{(8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8)}(T_{41},t)$ are eight augmentation coefficients for category 1, 2 and 3 $+(a_{44}^{''})^{(9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9)}(T_{45},t)$, $+(a_{46}^{''})^{(9,9,9)}(T_{45},t)$ are ninth augmentation coefficients for category 1, 2 and 3 | |
| $\frac{dT_{20}}{dt}=(b_{20})^{(3)}T_{21}-\begin{bmatrix}(b_{20}^{'})^{(3)}-(b_{20}^{''})^{(3)}(G_{23},t) & -(b_{16}^{''})^{(2,2,2)}(G_{19},t) & -(b_{13}^{''})^{(1,1,1,)}(G,t)\\ -(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t) & -(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t) & -(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t)\\ -(b_{36}^{''})^{(7,7,7,7)}(G_{39},t) & -(b_{40}^{''})^{(8,8,8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9,9)}(G_{47},t)\end{bmatrix}T_{20}$ | *70* |
| $\frac{dT_{21}}{dt}=(b_{21})^{(3)}T_{20}-\begin{bmatrix}(b_{21}^{'})^{(3)}-(b_{21}^{''})^{(3)}(G_{23},t) & -(b_{17}^{''})^{(2,2,2)}(G_{19},t) & -(b_{14}^{''})^{(1,1,1,)}(G,t)\\ -(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t) & -(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t) & -(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t)\\ -(b_{37}^{''})^{(7,7,7,7)}(G_{39},t) & -(b_{41}^{''})^{(8,8,8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9,9)}(G_{47},t)\end{bmatrix}T_{21}$ | *71* |
| $\frac{dT_{22}}{dt}=(b_{22})^{(3)}T_{21}-\begin{bmatrix}(b_{22}^{'})^{(3)}-(b_{22}^{''})^{(3)}(G_{23},t) & -(b_{18}^{''})^{(2,2,2)}(G_{19},t) & -(b_{15}^{''})^{(1,1,1,)}(G,t)\\ -(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t) & -(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t) & -(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t)\\ -(b_{38}^{''})^{(7,7,7,7)}(G_{39},t) & -(b_{42}^{''})^{(8,8,8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9)}(G_{47},t)\end{bmatrix}T_{22}$ | *72* |
| $-(b_{20}^{''})^{(3)}(G_{23},t)$, $-(b_{21}^{''})^{(3)}(G_{23},t)$, $-(b_{22}^{''})^{(3)}(G_{23},t)$ *are first detrition coefficients for category 1, 2 and 3* $-(b_{16}^{''})^{(2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2)}(G_{19},t)$ *are second detrition coefficients for category 1, 2 and 3* $-(b_{13}^{''})^{(1,1,1,)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,)}(G,t)$, $-(b_{15}^{''})^{(1,1,1,)}(G,t)$ *are third detrition coefficients for category 1,2 and 3* $-(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$ *are fourth detrition coefficients for category 1, 2 and 3* $-(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$ *are fifth detrition coefficients for category 1, 2 and 3* $-(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$ *are sixth detrition coefficients for category 1, 2 and 3* $-(b_{36}^{''})^{(7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3 $-(b_{40}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8)}(G_{43},t)$ are eight detrition coefficients for category 1, 2 and 3 $-(b_{46}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3 | |

$$\frac{dG_{24}}{dt} = (a_{24})^{(4)}G_{25} - \begin{bmatrix} (a'_{24})^{(4)} & +(a''_{24})^{(4)}(T_{25},t) & +(a''_{28})^{(5,5,)}(T_{29},t) & +(a''_{32})^{(6,6,)}(T_{33},t) \\ & +(a''_{13})^{(1,1,1,1)}(T_{14},t) & +(a''_{16})^{(2,2,2,2)}(T_{17},t) & +(a''_{20})^{(3,3,3,3)}(T_{21},t) \\ & +(a''_{36})^{(7,7,7,7,7)}(T_{37},t) & +(a''_{40})^{(8,8,8,8,8)}(T_{41},t) & +(a''_{44})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{24} \quad (73)$$

$$\frac{dG_{25}}{dt} = (a_{25})^{(4)}G_{24} - \begin{bmatrix} (a'_{25})^{(4)} & +(a''_{25})^{(4)}(T_{25},t) & +(a''_{29})^{(5,5,)}(T_{29},t) & +(a''_{33})^{(6,6)}(T_{33},t) \\ & +(a''_{14})^{(1,1,1,1)}(T_{14},t) & +(a''_{17})^{(2,2,2,2)}(T_{17},t) & +(a''_{21})^{(3,3,3,3)}(T_{21},t) \\ & +(a''_{37})^{(7,7,7,7,7)}(T_{37},t) & +(a''_{41})^{(8,8,8,8,8)}(T_{41},t) & +(a''_{45})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{25} \quad (74)$$

$$\frac{dG_{26}}{dt} = (a_{26})^{(4)}G_{25} - \begin{bmatrix} (a'_{26})^{(4)} & +(a''_{26})^{(4)}(T_{25},t) & +(a''_{30})^{(5,5,)}(T_{29},t) & +(a''_{34})^{(6,6,)}(T_{33},t) \\ & +(a''_{15})^{(1,1,1,1)}(T_{14},t) & +(a''_{18})^{(2,2,2,2)}(T_{17},t) & +(a''_{22})^{(3,3,3,3)}(T_{21},t) \\ & +(a''_{38})^{(7,7,7,7,7)}(T_{37},t) & +(a''_{42})^{(8,8,8,8,8)}(T_{41},t) & +(a''_{46})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{26} \quad (75)$$

$(a''_{24})^{(4)}(T_{25},t)$, $(a''_{25})^{(4)}(T_{25},t)$, $(a''_{26})^{(4)}(T_{25},t)$ *are first augmentation coefficients category* 1, 2 3

$+(a''_{28})^{(5,5,)}(T_{29},t)$, $+(a''_{29})^{(5,5,)}(T_{29},t)$, $+(a''_{30})^{(5,5,)}(T_{29},t)$ *are second augmentation coefficient for category* 1, 2 *and* 3

$+(a''_{32})^{(6,6,)}(T_{33},t)$, $+(a''_{33})^{(6,6,)}(T_{33},t)$, $+(a''_{34})^{(6,6,)}(T_{33},t)$ *are third augmentation coefficient for category* 1, 2 *and* 3

$+(a''_{13})^{(1,1,1,1)}(T_{14},t)$, $+(a''_{14})^{(1,1,1,1)}(T_{14},t)$, $+(a''_{15})^{(1,1,1,1)}(T_{14},t)$ *are fourth augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{16})^{(2,2,2,2)}(T_{17},t)$, $+(a''_{17})^{(2,2,2,2)}(T_{17},t)$, $+(a''_{18})^{(2,2,2,2)}(T_{17},t)$ *are fifth augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{20})^{(3,3,3,3)}(T_{21},t)$, $+(a''_{21})^{(3,3,3,3)}(T_{21},t)$, $+(a''_{22})^{(3,3,3,3)}(T_{21},t)$ *are sixth augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{36})^{(7,7,7,7,7)}(T_{37},t)$, $+(a''_{37})^{(7,7,7,7,7)}(T_{37},t)$, $+(a''_{38})^{(7,7,7,7,7)}(T_{37},t)$ *are seventh augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{40})^{(8,8,8,8,8)}(T_{41},t)$, $+(a''_{41})^{(8,8,8,8,8)}(T_{41},t)$, $+(a''_{42})^{(8,8,8,8,8)}(T_{41},t)$ *are eighth augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{46})^{(9,9,9,9)}(T_{45},t)$, $+(a''_{45})^{(9,9,9,9)}(T_{45},t)$, $+(a''_{44})^{(9,9,9,9)}(T_{45},t)$ are ninth detrition coefficients for category 1 2 3

$$\frac{dT_{24}}{dt} = (b_{24})^{(4)}T_{25} - \begin{bmatrix} (b'_{24})^{(4)} & -(b''_{24})^{(4)}(G_{27},t) & -(b''_{28})^{(5,5,)}(G_{31},t) & -(b''_{32})^{(6,6,)}(G_{35},t) \\ & -(b''_{13})^{(1,1,1,1)}(G,t) & -(b''_{16})^{(2,2,2,2)}(G_{19},t) & -(b''_{20})^{(3,3,3,3)}(G_{23},t) \\ & -(b''_{36})^{(7,7,7,7,7)}(G_{39},t) & -(b''_{40})^{(8,8,8,8,8)}(G_{43},t) & -(b''_{44})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{24} \quad (76)$$

$$\frac{dT_{25}}{dt} = (b_{25})^{(4)}T_{24} - \begin{bmatrix} (b'_{25})^{(4)} & -(b''_{25})^{(4)}(G_{27},t) & -(b''_{29})^{(5,5,)}(G_{31},t) & -(b''_{33})^{(6,6,)}(G_{35},t) \\ & -(b''_{14})^{(1,1,1,1)}(G,t) & -(b''_{17})^{(2,2,2,2)}(G_{19},t) & -(b''_{21})^{(3,3,3,3)}(G_{23},t) \\ & -(b''_{37})^{(7,7,7,7,7)}(G_{39},t) & -(b''_{41})^{(8,8,8,8,8)}(G_{43},t) & -(b''_{45})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{25} \quad (77)$$

$$\frac{dT_{26}}{dt} = (b_{26})^{(4)}T_{25} - \begin{bmatrix} (b'_{26})^{(4)}\boxed{-(b''_{26})^{(4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,)}(G_{35},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1)}(G,t)} & \boxed{-(b''_{18})^{(2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{38})^{(7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{26} \quad (78)$$

Where $\boxed{-(b''_{24})^{(4)}(G_{27},t)}$, $\boxed{-(b''_{25})^{(4)}(G_{27},t)}$, $\boxed{-(b''_{26})^{(4)}(G_{27},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{28})^{(5,5,)}(G_{31},t)}$, $\boxed{-(b''_{29})^{(5,5,)}(G_{31},t)}$, $\boxed{-(b''_{30})^{(5,5,)}(G_{31},t)}$ are second detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{32})^{(6,6,)}(G_{35},t)}$, $\boxed{-(b''_{33})^{(6,6,)}(G_{35},t)}$, $\boxed{-(b''_{34})^{(6,6,)}(G_{35},t)}$ are third detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{13})^{(1,1,1,1)}(G,t)}$, $\boxed{-(b''_{14})^{(1,1,1,1)}(G,t)}$, $\boxed{-(b''_{15})^{(1,1,1,1)}(G,t)}$ are fourth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{16})^{(2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{17})^{(2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{18})^{(2,2,2,2)}(G_{19},t)}$ are fifth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{20})^{(3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{21})^{(3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{22})^{(3,3,3,3)}(G_{23},t)}$ are sixth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{36})^{(7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{37})^{(7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{38})^{(7,7,7,7,7)}(G_{39},t)}$ are seventh detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{40})^{(8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{41})^{(8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{42})^{(8,8,8,8,8)}(G_{43},t)}$ are eighth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{46})^{(9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{45})^{(9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{44})^{(9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1 2 3

$$\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \begin{bmatrix} (a'_{28})^{(5)}\boxed{+(a''_{28})^{(5)}(T_{29},t)} & \boxed{+(a''_{24})^{(4,4,)}(T_{25},t)} & \boxed{+(a''_{32})^{(6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{36})^{(7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{40})^{(8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{44})^{(9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{28} \quad (79)$$

$$\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \begin{bmatrix} (a'_{29})^{(5)}\boxed{+(a''_{29})^{(5)}(T_{29},t)} & \boxed{+(a''_{25})^{(4,4,)}(T_{25},t)} & \boxed{+(a''_{33})^{(6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{14})^{(1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{17})^{(2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{21})^{(3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{37})^{(7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{41})^{(8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{45})^{(9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{29} \quad (80)$$

$$\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \begin{bmatrix} (a'_{30})^{(5)}\boxed{+(a''_{30})^{(5)}(T_{29},t)} & \boxed{+(a''_{26})^{(4,4,)}(T_{25},t)} & \boxed{+(a''_{34})^{(6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{15})^{(1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{18})^{(2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{22})^{(3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{38})^{(7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{42})^{(8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{46})^{(9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{30} \quad (81)$$

| | |
|---|---|
| *Where* $\boxed{+(a_{28}^{''})^{(5)}(T_{29},t)}$, $\boxed{+(a_{29}^{''})^{(5)}(T_{29},t)}$, $\boxed{+(a_{30}^{''})^{(5)}(T_{29},t)}$ *are first augmentation coefficients for category* 1, 2 *and* 3<br>*And* $\boxed{+(a_{24}^{''})^{(4,4,)}(T_{25},t)}$, $\boxed{+(a_{25}^{''})^{(4,4,)}(T_{25},t)}$, $\boxed{+(a_{26}^{''})^{(4,4,)}(T_{25},t)}$ *are second augmentation coefficient for category* 1, 2 *and* 3<br>$\boxed{+(a_{32}^{''})^{(6,6,6)}(T_{33},t)}$, $\boxed{+(a_{33}^{''})^{(6,6,6)}(T_{33},t)}$, $\boxed{+(a_{34}^{''})^{(6,6,6)}(T_{33},t)}$ *are third augmentation coefficient for category* 1, 2 *and* 3<br>$\boxed{+(a_{13}^{''})^{(1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a_{14}^{''})^{(1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a_{15}^{''})^{(1,1,1,1,1)}(T_{14},t)}$ are fourth augmentation coefficients for category 1,2, and 3<br>$\boxed{+(a_{16}^{''})^{(2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a_{17}^{''})^{(2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a_{18}^{''})^{(2,2,2,2,2)}(T_{17},t)}$ are fifth augmentation coefficients for category 1,2,and 3<br>$\boxed{+(a_{20}^{''})^{(3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a_{21}^{''})^{(3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a_{22}^{''})^{(3,3,3,3,3)}(T_{21},t)}$ are sixth augmentation coefficients for category 1,2, 3<br>$\boxed{+(a_{36}^{''})^{(7,7,7,7,7,7)}(T_{37},t)}$, $\boxed{+(a_{37}^{''})^{(7,7,7,7,7,7)}(T_{37},t)}$, $\boxed{+(a_{38}^{''})^{(7,7,7,7,7,7)}(T_{37},t)}$ are seventh augmentation coefficients for category 1,2, 3<br>$\boxed{+(a_{40}^{''})^{(8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{41}^{''})^{(8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{42}^{''})^{(8,8,8,8,8,8)}(T_{41},t)}$ are eighth augmentation coefficients for category 1,2, 3<br>$\boxed{+(a_{46}^{''})^{(9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{45}^{''})^{(9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{44}^{''})^{(9,9,9,9,9)}(T_{45},t)}$ are ninth augmentation coefficients for category 1,2, 3 | |
| $\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \left[\begin{matrix} (b_{28}^{'})^{(5)}\boxed{-(b_{28}^{''})^{(5)}(G_{31},t)} & \boxed{-(b_{24}^{''})^{(4,4,)}(G_{27},t)} & \boxed{-(b_{32}^{''})^{(6,6,6)}(G_{35},t)} \\ \boxed{-(b_{13}^{''})^{(1,1,1,1,1)}(G,t)} & \boxed{-(b_{16}^{''})^{(2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{20}^{''})^{(3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{36}^{''})^{(7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{40}^{''})^{(8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b_{44}^{''})^{(9,9,9,9,9)}(G_{47},t)} \end{matrix}\right] T_{28}$ | *82* |
| $\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \left[\begin{matrix} (b_{29}^{'})^{(5)}\boxed{-(b_{29}^{''})^{(5)}(G_{31},t)} & \boxed{-(b_{25}^{''})^{(4,4,)}(G_{27},t)} & \boxed{-(b_{33}^{''})^{(6,6,6)}(G_{35},t)} \\ \boxed{-(b_{14}^{''})^{(1,1,1,1,1)}(G,t)} & \boxed{-(b_{17}^{''})^{(2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{21}^{''})^{(3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{37}^{''})^{(7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{41}^{''})^{(8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b_{45}^{''})^{(9,9,9,9,9)}(G_{47},t)} \end{matrix}\right] T_{29}$ | *83* |
| $\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \left[\begin{matrix} (b_{30}^{'})^{(5)}\boxed{-(b_{30}^{''})^{(5)}(G_{31},t)} & \boxed{-(b_{26}^{''})^{(4,4,)}(G_{27},t)} & \boxed{-(b_{34}^{''})^{(6,6,6)}(G_{35},t)} \\ \boxed{-(b_{15}^{''})^{(1,1,1,1,1,)}(G,t)} & \boxed{-(b_{18}^{''})^{(2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{22}^{''})^{(3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{38}^{''})^{(7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{42}^{''})^{(8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b_{46}^{''})^{(9,9,9,9,9)}(G_{47},t)} \end{matrix}\right] T_{30}$ | *84* |
| *where* $\boxed{-(b_{28}^{''})^{(5)}(G_{31},t)}$, $\boxed{-(b_{29}^{''})^{(5)}(G_{31},t)}$, $\boxed{-(b_{30}^{''})^{(5)}(G_{31},t)}$ *are first detrition coefficients for category* 1, 2 *and* 3<br>$\boxed{-(b_{24}^{''})^{(4,4,)}(G_{27},t)}$, $\boxed{-(b_{25}^{''})^{(4,4,)}(G_{27},t)}$, $\boxed{-(b_{26}^{''})^{(4,4,)}(G_{27},t)}$ *are second detrition coefficients for category* 1,2 *and* 3<br>$\boxed{-(b_{32}^{''})^{(6,6,6)}(G_{35},t)}$, $\boxed{-(b_{33}^{''})^{(6,6,6)}(G_{35},t)}$, $\boxed{-(b_{34}^{''})^{(6,6,6)}(G_{35},t)}$ *are third detrition coefficients for category* 1,2 *and* 3<br>$\boxed{-(b_{13}^{''})^{(1,1,1,1,1)}(G,t)}$, $\boxed{-(b_{14}^{''})^{(1,1,1,1,1)}(G,t)}$, $\boxed{-(b_{15}^{''})^{(1,1,1,1,1,)}(G,t)}$ are fourth detrition coefficients for category 1,2, and 3<br>$\boxed{-(b_{16}^{''})^{(2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b_{17}^{''})^{(2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b_{18}^{''})^{(2,2,2,2,2)}(G_{19},t)}$ are fifth detrition coefficients | |

| | |
|---|---|
| for category 1,2, and 3<br>$-(b_{20}^{''})^{(3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3)}(G_{23},t)$ are sixth detrition coefficients for category 1,2, and 3<br>$-(b_{36}^{''})^{(7,7,7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1,2, and 3<br>$-(b_{42}^{''})^{(8,8,8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8,8,8)}(G_{43},t)$, $-(b_{40}^{''})^{(8,8,8,8,8,8)}(G_{43},t)$ are eighth detrition coefficients for category 1,2, and 3<br>$-(b_{46}^{''})^{(9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1,2, and 3 | |
| $\frac{dG_{32}}{dt} = (a_{32})^{(6)}G_{33} - \begin{bmatrix} (a_{32}^{'})^{(6)}+(a_{32}^{''})^{(6)}(T_{33},t) & +(a_{28}^{''})^{(5,5,5)}(T_{29},t) & +(a_{24}^{''})^{(4,4,4,)}(T_{25},t) \\ +(a_{13}^{''})^{(1,1,1,1,1,1)}(T_{14},t) & +(a_{16}^{''})^{(2,2,2,2,2,2)}(T_{17},t) & +(a_{20}^{''})^{(3,3,3,3,3,3)}(T_{21},t) \\ +(a_{36}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t) & +(a_{40}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a_{44}^{''})^{(9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{32}$ | *85* |
| $\frac{dG_{33}}{dt} = (a_{33})^{(6)}G_{32} - \begin{bmatrix} (a_{33}^{'})^{(6)}+(a_{33}^{''})^{(6)}(T_{33},t) & +(a_{29}^{''})^{(5,5,5)}(T_{29},t) & +(a_{25}^{''})^{(4,4,4,)}(T_{25},t) \\ +(a_{14}^{''})^{(1,1,1,1,1,1)}(T_{14},t) & +(a_{17}^{''})^{(2,2,2,2,2,2)}(T_{17},t) & +(a_{21}^{''})^{(3,3,3,3,3,3)}(T_{21},t) \\ +(a_{37}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t) & +(a_{41}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a_{45}^{''})^{(9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{33}$ | *86* |
| $\frac{dG_{34}}{dt} = (a_{34})^{(6)}G_{33} - \begin{bmatrix} (a_{34}^{'})^{(6)}+(a_{34}^{''})^{(6)}(T_{33},t) & +(a_{30}^{''})^{(5,5,5)}(T_{29},t) & +(a_{26}^{''})^{(4,4,4,)}(T_{25},t) \\ +(a_{15}^{''})^{(1,1,1,1,1,1)}(T_{14},t) & +(a_{18}^{''})^{(2,2,2,2,2,2)}(T_{17},t) & +(a_{22}^{''})^{(3,3,3,3,3,3)}(T_{21},t) \\ +(a_{38}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t) & +(a_{42}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a_{46}^{''})^{(9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{34}$ | *87* |
| $+(a_{32}^{''})^{(6)}(T_{33},t)$, $+(a_{33}^{''})^{(6)}(T_{33},t)$, $+(a_{34}^{''})^{(6)}(T_{33},t)$ *are first augmentation coefficients for category* 1, 2 *and* 3<br>$+(a_{28}^{''})^{(5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5)}(T_{29},t)$ *are second augmentation coefficients for category* 1, 2 *and* 3<br>$+(a_{24}^{''})^{(4,4,4,)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,)}(T_{25},t)$ *are third augmentation coefficients for category* 1, 2 *and* 3<br>$+(a_{13}^{''})^{(1,1,1,1,1,1)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1,1,1)}(T_{14},t)$ - are fourth augmentation coefficients<br>$+(a_{16}^{''})^{(2,2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2,2)}(T_{17},t)$ - fifth augmentation coefficients<br>$+(a_{20}^{''})^{(3,3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3,3)}(T_{21},t)$ sixth augmentation coefficients<br>$+(a_{36}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)$, $+(a_{37}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)$ seventh augmentation coefficients<br>$+(a_{40}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)$ Eighth augmentation coefficients<br>$+(a_{44}^{''})^{(9,9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9,9)}(T_{45},t)$, $+(a_{46}^{''})^{(9,9,9,9,9,9)}(T_{45},t)$ ninth augmentation | |

coefficients

$$\frac{dT_{32}}{dt} = (b_{32})^{(6)}T_{33} - \begin{bmatrix} (b'_{32})^{(6)}\boxed{-(b''_{32})^{(6)}(G_{35},t)} & \boxed{-(b''_{28})^{(5,5,5)}(G_{31},t)} & \boxed{-(b''_{24})^{(4,4,4,)}(G_{27},t)} \\ \boxed{-(b''_{13})^{(1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{16})^{(2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{20})^{(3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{36})^{(7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{40})^{(8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{44})^{(9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{32} \quad (88)$$

$$\frac{dT_{33}}{dt} = (b_{33})^{(6)}T_{32} - \begin{bmatrix} (b'_{33})^{(6)}\boxed{-(b''_{33})^{(6)}(G_{35},t)} & \boxed{-(b''_{29})^{(5,5,5)}(G_{31},t)} & \boxed{-(b''_{25})^{(4,4,4,)}(G_{27},t)} \\ \boxed{-(b''_{14})^{(1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{17})^{(2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{21})^{(3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{37})^{(7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{41})^{(8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{45})^{(9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{33} \quad (89)$$

$$\frac{dT_{34}}{dt} = (b_{34})^{(6)}T_{33} - \begin{bmatrix} (b'_{34})^{(6)}\boxed{-(b''_{34})^{(6)}(G_{35},t)} & \boxed{-(b''_{30})^{(5,5,5)}(G_{31},t)} & \boxed{-(b''_{26})^{(4,4,4,)}(G_{27},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{18})^{(2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{38})^{(7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{34} \quad (90)$$

$\boxed{-(b''_{32})^{(6)}(G_{35},t)}, \boxed{-(b''_{33})^{(6)}(G_{35},t)}, \boxed{-(b''_{34})^{(6)}(G_{35},t)}$ *are first detrition coefficients for category* 1, 2 *and* 3

$\boxed{-(b''_{28})^{(5,5,5)}(G_{31},t)}, \boxed{-(b''_{29})^{(5,5,5)}(G_{31},t)}, \boxed{-(b''_{30})^{(5,5,5)}(G_{31},t)}$ *are second detrition coefficients for category* 1, 2 *and* 3

$\boxed{-(b''_{24})^{(4,4,4,)}(G_{27},t)}, \boxed{-(b''_{25})^{(4,4,4,)}(G_{27},t)}, \boxed{-(b''_{26})^{(4,4,4,)}(G_{27},t)}$ *are third detrition coefficients for category* 1,2 *and* 3

$\boxed{-(b''_{13})^{(1,1,1,1,1,1)}(G,t)}, \boxed{-(b''_{14})^{(1,1,1,1,1,1)}(G,t)}, \boxed{-(b''_{15})^{(1,1,1,1,1,1)}(G,t)}$ are fourth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{16})^{(2,2,2,2,2,2)}(G_{19},t)}, \boxed{-(b''_{17})^{(2,2,2,2,2,2)}(G_{19},t)}, \boxed{-(b''_{18})^{(2,2,2,2,2,2)}(G_{19},t)}$ are fifth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{20})^{(3,3,3,3,3,3)}(G_{23},t)}, \boxed{-(b''_{21})^{(3,3,3,3,3,3)}(G_{23},t)}, \boxed{-(b''_{22})^{(3,3,3,3,3,3)}(G_{23},t)}$ are sixth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{36})^{(7,7,7,7,7,7,7)}(G_{39},t)}, \boxed{-(b''_{37})^{(7,7,7,7,7,7,7)}(G_{39},t)}, \boxed{-(b''_{38})^{(7,7,7,7,7,7,7)}(G_{39},t)}$ are seventh detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{40})^{(8,8,8,8,8,8,8)}(G_{43},t)}, \boxed{-(b''_{41})^{(8,8,8,8,8,8,8)}(G_{43},t)}, \boxed{-(b''_{42})^{(8,8,8,8,8,8,8)}(G_{43},t)}$ are eighth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{46})^{(9,9,9,9,9,9)}(G_{47},t)}, \boxed{-(b''_{45})^{(9,9,9,9,9,9)}(G_{47},t)}, \boxed{-(b''_{44})^{(9,9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2, and 3

$$\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \begin{bmatrix} (a'_{36})^{(7)}\boxed{+(a''_{36})^{(7)}(T_{37},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{24})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8,)}(T_{41},t)} & \boxed{+(a''_{44})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{13} \quad (91)$$

$$\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \begin{bmatrix} (a_{37}')^{(7)}\boxed{+(a_{37}'')^{(7)}(T_{37},t)} & \boxed{+(a_{17}'')^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{21}'')^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{25}'')^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{29}'')^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{33}'')^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{13}'')^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{41}'')^{(8,8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{45}'')^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{14} \quad (92)$$

$$\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \begin{bmatrix} (a_{38}')^{(7)}\boxed{+(a_{38}'')^{(7)}(T_{37},t)} & \boxed{+(a_{18}'')^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{22}'')^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{26}'')^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{30}'')^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{34}'')^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{15}'')^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{42}'')^{(8,8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{46}'')^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{15} \quad (93)$$

Where $\boxed{(a_{36}'')^{(7)}(T_{37},t)}$, $\boxed{(a_{37}'')^{(7)}(T_{37},t)}$, $\boxed{(a_{38}'')^{(7)}(T_{37},t)}$ are first augmentation coefficients for category 1, 2 and 3

$\boxed{+(a_{16}'')^{(2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a_{17}'')^{(2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a_{18}'')^{(2,2,2,2,2,2,2)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{20}'')^{(3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a_{21}'')^{(3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a_{22}'')^{(3,3,3,3,3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{24}'')^{(4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{25}'')^{(4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{26}'')^{(4,4,4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{28}'')^{(5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{29}'')^{(5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{30}'')^{(5,5,5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{32}'')^{(6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{33}'')^{(6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{34}'')^{(6,6,6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{13}'')^{(1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a_{13}'')^{(1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a_{15}'')^{(1,1,1,1,1,1,1)}(T_{14},t)}$ are seventh augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{42}'')^{(8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{41}'')^{(8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{40}'')^{(8,8,8,8,8,8,8,8,)}(T_{41},t)}$ are eighth augmentation coefficient for 1,2,3

$\boxed{+(a_{46}'')^{(9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{45}'')^{(9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{44}'')^{(9,9,9,9,9,9,9)}(T_{45},t)}$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{36}}{dt} = (b_{36})^{(7)}T_{37} - \begin{bmatrix} (b_{36}')^{(7)}\boxed{-(b_{36}'')^{(7)}(G_{39},t)} & \boxed{-(b_{16}'')^{(2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{20}'')^{(3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{24}'')^{(4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{28}'')^{(5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{32}'')^{(6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{13}'')^{(1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{40}'')^{(8,8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b_{44}'')^{(9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{13} \quad (94)$$

$$\frac{dT_{37}}{dt} =$$

$$(b_{37})^{(7)}T_{36} - \begin{bmatrix} (b'_{37})^{(7)}\boxed{-(b''_{37})^{(7)}(G_{39},t)} & \boxed{-(b''_{17})^{(2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{21})^{(3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{25})^{(4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{29})^{(5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{33})^{(6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{14})^{(1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{41})^{(8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{45})^{(9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{14}$$

$$\frac{dT_{38}}{dt} =$$

$$(b_{38})^{(7)}T_{37} - \begin{bmatrix} (b'_{38})^{(7)}\boxed{-(b''_{38})^{(7)}(G_{39},t)} & \boxed{-(b''_{18})^{(2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{26})^{(4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{42})^{(8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{15}$$

Where $\boxed{-(b''_{36})^{(7)}(G_{39},t)}$, $\boxed{-(b''_{37})^{(7)}(G_{39},t)}$, $\boxed{-(b''_{38})^{(7)}(G_{39},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{16})^{(2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{17})^{(2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{18})^{(2,2,2,2,2,2,2)}(G_{19},t)}$ are second detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{20})^{(3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{21})^{(3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{22})^{(3,3,3,3,3,3,3)}(G_{23},t)}$ are third detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{24})^{(4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{25})^{(4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{26})^{(4,4,4,4,4,4,4)}(G_{27},t)}$ are fourth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{28})^{(5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{29})^{(5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{30})^{(5,5,5,5,5,5,5)}(G_{31},t)}$ are fifth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{32})^{(6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b''_{33})^{(6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b''_{34})^{(6,6,6,6,6,6,6)}(G_{35},t)}$ are sixth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{15})^{(1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{14})^{(1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{13})^{(1,1,1,1,1,1,1)}(G,t)}$ are seventh detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{40})^{(8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{41})^{(8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{42})^{(8,8,8,8,8,8,8)}(G_{43},t)}$ are eighth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{46})^{(9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{45})^{(9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{44})^{(9,9,9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2 and 3

$$\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \begin{bmatrix} (a'_{40})^{(8)}\boxed{+(a''_{40})^{(8)}(T_{41},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{24})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{36})^{(7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{44})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{13} \quad (95)$$

$$\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \begin{bmatrix} (a'_{41})^{(8)}\boxed{+(a''_{41})^{(8)}(T_{41},t)} & \boxed{+(a''_{17})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{21})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{25})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{29})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{33})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{37})^{(7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{45})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{14}$$

$$\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \begin{bmatrix} (a'_{42})^{(8)}\boxed{+(a''_{42})^{(8)}(T_{41},t)} & \boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{46})^{(9,9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{15}$$

Where $\boxed{+(a''_{40})^{(8)}(T_{41},t)}$, $\boxed{+(a''_{41})^{(8)}(T_{41},t)}$, $\boxed{+(a''_{42})^{(8)}(T_{41},t)}$ are first augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{16})^{(2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{17})^{(2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{20})^{(3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{21})^{(3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{24})^{(4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{25})^{(4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{28})^{(5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{29})^{(5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{32})^{(6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{33})^{(6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{13})^{(1,1,1,1,1,1,1,1)}(T_{14},t)}$ $\boxed{+(a''_{14})^{(1,1,1,1,1,1,1,1)}(T_{14},t)}$ $\boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1)}(T_{14},t)}$ are seventh augmentation coefficient for 1,2,3

$\boxed{+(a''_{36})^{(7,7,7,7,7,7,7,7)}(T_{37},t)}$, $\boxed{+(a''_{37})^{(7,7,7,7,7,7,7,7)}(T_{37},t)}$, $\boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7)}(T_{37},t)}$ are eighth augmentation coefficient for 1,2,3

$\boxed{+(a''_{46})^{(9,9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a''_{45})^{(9,9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a''_{44})^{(9,9,9,9,9,9,9,9)}(T_{45},t)}$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{40}}{dt} = (b_{40})^{(8)}T_{41} - \begin{bmatrix} (b'_{40})^{(8)}\boxed{-(b''_{40})^{(8)}(G_{43},t)} & \boxed{-(b''_{16})^{(2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{20})^{(3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{24})^{(4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{28})^{(5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{32})^{(6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{13})^{(1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{36})^{(7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{44})^{(9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{13}$$

$$\frac{dT_{41}}{dt} = (b_{41})^{(8)}T_{40} - \begin{bmatrix} (b'_{41})^{(8)}\boxed{-(b''_{41})^{(8)}(G_{43},t)} & \boxed{-(b''_{17})^{(2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{21})^{(3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{25})^{(4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{29})^{(5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{33})^{(6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{14})^{(1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{37})^{(7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{45})^{(9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{14}$$

$$\frac{dT_{42}}{dt} =$$

$$(b_{42})^{(8)}T_{41} - \begin{bmatrix} (b'_{42})^{(8)}\boxed{-(b''_{42})^{(8)}(G_{43},t)} & \boxed{-(b''_{18})^{(2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{26})^{(4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{38})^{(7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{46})^{(9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{15}$$

Where $\boxed{-(b''_{36})^{(7)}(G_{39},t)}$, $\boxed{-(b''_{37})^{(7)}(G_{39},t)}$, $\boxed{-(b''_{38})^{(7)}(G_{39},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{16})^{(2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{17})^{(2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{18})^{(2,2,2,2,2,2,2,2)}(G_{19},t)}$ are second detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{20})^{(3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{21})^{(3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{22})^{(3,3,3,3,3,3,3,3)}(G_{23},t)}$ are third detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{24})^{(4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{25})^{(4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{26})^{(4,4,4,4,4,4,4,4)}(G_{27},t)}$ are fourth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{28})^{(5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{29})^{(5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{30})^{(5,5,5,5,5,5,5,5)}(G_{31},t)}$ are fifth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{32})^{(6,6,6,6,)}(G_{35},t)}$, $\boxed{-(b''_{33})^{(6,6,6,6,)}(G_{35},t)}$, $\boxed{-(b''_{15})^{(1,1,1,1,1,1,1,1)}(G,t)}$ are sixth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{13})^{(1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{14})^{(1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{38})^{(7,7,)}(G_{39},t)}$ are seventh detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{36})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{37})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{38})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$ are eighth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{44})^{(9,9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{45})^{(9,9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{46})^{(9,9,9,9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2 and 3

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$$\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \begin{bmatrix} (a'_{44})^{(9)}\boxed{+(a''_{44})^{(9)}(T_{45},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{24})^{(4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{36})^{(7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8)}(T_{41},t)} \end{bmatrix} G_{13}$$

$$\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \begin{bmatrix} (a'_{45})^{(9)}\boxed{+(a''_{45})^{(9)}(T_{45},t)} & \boxed{+(a''_{17})^{(2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{21})^{(3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{25})^{(4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{29})^{(5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{33})^{(6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{14})^{(1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{37})^{(7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{41})^{(8,8,8,8,8,8,8,8)}(T_{41},t)} \end{bmatrix} G_{14}$$

$$\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \begin{bmatrix} (a'_{46})^{(9)}\boxed{+(a''_{46})^{(9)}(T_{37},t)} & \boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{42})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)} \end{bmatrix} G_{15}$$

Where $\boxed{+(a''_{44})^{(9)}(T_{45},t)}$, $\boxed{+(a''_{45})^{(9)}(T_{45},t)}$, $\boxed{+(a''_{46})^{(9)}(T_{37},t)}$ are first augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{16})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{17})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{20})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{21})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{24})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{25})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{28})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{29})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{32})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{33})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{13})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a''_{14})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)}$ are Seventh augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)}$ $\boxed{+(a''_{37})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)}$ $\boxed{+(a''_{36})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)}$ are eighth augmentation coefficient for 1,2,3

$\boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a''_{42})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a''_{41})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)}$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{44}}{dt} = (b_{44})^{(9)}T_{45} - \begin{bmatrix} (b'_{44})^{(9)}\boxed{-(b''_{44})^{(9)}(G_{47},t)} & \boxed{-(b''_{16})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{20})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{24})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{28})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{32})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{13})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{36})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{40})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{bmatrix} T_{13}$$

$$\frac{dT_{45}}{dt} = (b_{45})^{(9)}T_{44} - \begin{bmatrix} (b'_{45})^{(9)}\boxed{-(b''_{45})^{(9)}(G_{47},t)} & \boxed{-(b''_{17})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{21})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{25})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{29})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{33})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{14})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{37})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{41})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{bmatrix} T_{14}$$

$$\frac{dT_{46}}{dt} =$$

$$(b_{46})^{(9)}T_{45} - \begin{bmatrix} (b'_{46})^{(9)}\boxed{-(b''_{46})^{(9)}(G_{47},t)} & \boxed{-(b''_{18})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{26})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{38})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{bmatrix} T_{15}$$

Where $\boxed{-(b''_{44})^{(9)}(G_{47},t)}$, $\boxed{-(b''_{45})^{(9)}(G_{47},t)}$, $\boxed{-(b''_{46})^{(9)}(G_{47},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{16})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{17})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{18})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)}$ are second detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{20})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{21})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{22})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)}$ are third detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{24})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{25})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{26})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)}$ are fourth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{28})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{29})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{30})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)}$ are fifth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{32})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b''_{33})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b''_{34})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)}$ are sixth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{15})^{(1,1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{14})^{(1,1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{13})^{(1,1,1,1,1,1,1,1,1)}(G,t)}$ are seventh detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{37})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{36})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{38})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)}$ are eighth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{42})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{41})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{40})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$ are ninth detrition coefficients for category 1, 2 and 3

Where we suppose

$(a_i)^{(1)}, (a'_i)^{(1)}, (a''_i)^{(1)}, (b_i)^{(1)}, (b'_i)^{(1)}, (b''_i)^{(1)} > 0, \quad i,j = 13,14,15$ *97*

The functions $(a''_i)^{(1)}, (b''_i)^{(1)}$ are positive continuous increasing and bounded.

**Definition of** $(p_i)^{(1)}$, $(r_i)^{(1)}$:

$$(a''_i)^{(1)}(T_{14},t) \le (p_i)^{(1)} \le (\hat{A}_{13})^{(1)}$$

$$(b''_i)^{(1)}(G,t) \le (r_i)^{(1)} \le (b'_i)^{(1)} \le (\hat{B}_{13})^{(1)}$$

$lim_{T_2\to\infty}(a''_i)^{(1)}(T_{14},t) = (p_i)^{(1)}$ 98

$\lim_{G\to\infty}(b''_i)^{(1)}(G,t) = (r_i)^{(1)}$

**Definition of** $(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}$ :

Where $\boxed{(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}}$ are positive constants and $\boxed{i = 13,14,15}$

| | |
|---|---|
| They satisfy Lipschitz condition:<br>$\|(a_i'')^{(1)}(T_{14}', t) - (a_i'')^{(1)}(T_{14}, t)\| \leq (\hat{k}_{13})^{(1)}\|T_{14} - T_{14}'\|e^{-(\hat{M}_{13})^{(1)}t}$<br>$\|(b_i'')^{(1)}(G', t) - (b_i'')^{(1)}(G, t)\| < (\hat{k}_{13})^{(1)}\|G - G'\|e^{-(\hat{M}_{13})^{(1)}t}$ | 99 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(1)}(T_{14}', t)$ and $(a_i'')^{(1)}(T_{14}, t)$ . $(T_{14}', t)$ and $(T_{14}, t)$ are points belonging to the interval $[(\hat{k}_{13})^{(1)}, (\hat{M}_{13})^{(1)}]$. It is to be noted that $(a_i'')^{(1)}(T_{14}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{13})^{(1)} = 1$ then the function $(a_i'')^{(1)}(T_{14}, t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}$ :<br>$(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)},$ are positive constants<br>$\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}} < 1$ | 100 |
| **Definition of** $(\hat{P}_{13})^{(1)}, (\hat{Q}_{13})^{(1)}$ :<br>There exists two constants $(\hat{P}_{13})^{(1)}$ and $(\hat{Q}_{13})^{(1)}$ which together With $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}, (\hat{A}_{13})^{(1)}$ and $(\hat{B}_{13})^{(1)}$ and the constants $(a_i)^{(1)}, (a_i')^{(1)}, (b_i)^{(1)}, (b_i')^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}, i = 13,14,15,$ satisfy the inequalities<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(a_i)^{(1)} + (a_i')^{(1)} + (\hat{A}_{13})^{(1)} + (\hat{P}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(b_i)^{(1)} + (b_i')^{(1)} + (\hat{B}_{13})^{(1)} + (\hat{Q}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$ | 101 |
| Where we suppose | |
| $(a_i)^{(2)}, (a_i')^{(2)}, (a_i'')^{(2)}, (b_i)^{(2)}, (b_i')^{(2)}, (b_i'')^{(2)} > 0, \quad i, j = 16,17,18$ | |
| The functions $(a_i'')^{(2)}, (b_i'')^{(2)}$ are positive continuous increasing and bounded. | |
| **Definition of** $(p_i)^{(2)}$, $(r_i)^{(2)}$: | |
| $(a_i'')^{(2)}(T_{17}, t) \leq (p_i)^{(2)} \leq (\hat{A}_{16})^{(2)}$ | 102 |
| $(b_i'')^{(2)}(G_{19}, t) \leq (r_i)^{(2)} \leq (b_i')^{(2)} \leq (\hat{B}_{16})^{(2)}$ | 103 |
| $\lim_{T_2 \to \infty}(a_i'')^{(2)}(T_{17}, t) = (p_i)^{(2)}$ | 104 |
| $\lim_{G \to \infty}(b_i'')^{(2)}((G_{19}), t) = (r_i)^{(2)}$ | 105 |
| **Definition of** $(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}$ : | 106 |

| | |
|---|---|
| Where $\boxed{(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}, (p_i)^{(2)}, (r_i)^{(2)}}$ are positive constants and $\boxed{i = 16,17,18}$ | |
| They satisfy Lipschitz condition: | |
| $\lvert(a_i^{''})^{(2)}(T_{17}^{'}, t) - (a_i^{''})^{(2)}(T_{17}, t)\rvert \leq (\hat{k}_{16})^{(2)} \lvert T_{17} - T_{17}^{'}\rvert e^{-(\hat{M}_{16})^{(2)}t}$ | 107 |
| $\lvert(b_i^{''})^{(2)}((G_{19})^{'}, t) - (b_i^{''})^{(2)}((G_{19}), t)\rvert < (\hat{k}_{16})^{(2)} \lvert\lvert(G_{19}) - (G_{19})^{'}\rvert\rvert e^{-(\hat{M}_{16})^{(2)}t}$ | 108 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(2)}(T_{17}^{'}, t)$ and $(a_i^{''})^{(2)}(T_{17}, t)$ . $(T_{17}^{'}, t)$ and $(T_{17}, t)$ are points belonging to the interval $\left[(\hat{k}_{16})^{(2)}, (\hat{M}_{16})^{(2)}\right]$. It is to be noted that $(a_i^{''})^{(2)}(T_{17}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{16})^{(2)} = 1$ then the function $(a_i^{''})^{(2)}(T_{17}, t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$ : | |
| $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$, are positive constants<br><br>$\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}$ , $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ | 109 |
| **Definition of** $(\hat{P}_{13})^{(2)}, (\hat{Q}_{13})^{(2)}$ :<br><br>There exists two constants $(\hat{P}_{16})^{(2)}$ and $(\hat{Q}_{16})^{(2)}$ which together with $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}, (\hat{A}_{16})^{(2)} and (\hat{B}_{16})^{(2)}$ and the constants $(a_i)^{(2)}, (a_i^{'})^{(2)}, (b_i)^{(2)}, (b_i^{'})^{(2)}, (p_i)^{(2)}, (r_i)^{(2)}, i = 16,17,18,$<br><br>satisfy the inequalities | |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(a_i)^{(2)} + (a_i^{'})^{(2)} + (\hat{A}_{16})^{(2)} + (\hat{P}_{16})^{(2)} (\hat{k}_{16})^{(2)}] < 1$ | 110 |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(b_i)^{(2)} + (b_i^{'})^{(2)} + (\hat{B}_{16})^{(2)} + (\hat{Q}_{16})^{(2)} (\hat{k}_{16})^{(2)}] < 1$ | 111 |
| Where we suppose | |
| $(a_i)^{(3)}, (a_i^{'})^{(3)}, (a_i^{''})^{(3)}, (b_i)^{(3)}, (b_i^{'})^{(3)}, (b_i^{''})^{(3)} > 0, \quad i, j = 20,21,22$<br><br>The functions $(a_i^{''})^{(3)}, (b_i^{''})^{(3)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(3)}$, $(r_i)^{(3)}$:<br><br>$(a_i^{''})^{(3)}(T_{21}, t) \leq (p_i)^{(3)} \leq (\hat{A}_{20})^{(3)}$<br><br>$(b_i^{''})^{(3)}(G_{23}, t) \leq (r_i)^{(3)} \leq (b_i^{'})^{(3)} \leq (\hat{B}_{20})^{(3)}$ | 112 |
| $lim_{T_2 \to \infty}(a_i^{''})^{(3)}(T_{21}, t) = (p_i)^{(3)}$ | 113 |

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|---|---|
| $\lim_{G\to\infty}(b_i'')^{(3)}(G_{23},t) = (r_i)^{(3)}$<br><br>**Definition of** $(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}$ :<br><br>Where $\boxed{(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}}$ are positive constants and $\boxed{i = 20,21,22}$ | |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i'')^{(3)}(T_{21}',t) - (a_i'')^{(3)}(T_{21},t)\rvert \le (\hat{k}_{20})^{(3)}\lvert T_{21} - T_{21}'\rvert e^{-(\hat{M}_{20})^{(3)}t}$<br><br>$\lvert(b_i'')^{(3)}(G_{23}',t) - (b_i'')^{(3)}(G_{23},t)\rvert < (\hat{k}_{20})^{(3)}\lvert\lvert G_{23} - G_{23}'\rvert\rvert e^{-(\hat{M}_{20})^{(3)}t}$ | 114 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(3)}(T_{21}',t)$ and$(a_i'')^{(3)}(T_{21},t)$ . $(T_{21}',t)$ And $(T_{21},t)$ are points belonging to the interval $\left[(\hat{k}_{20})^{(3)}, (\hat{M}_{20})^{(3)}\right]$. It is to be noted that $(a_i'')^{(3)}(T_{21},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{20})^{(3)} = 1$ then the function $(a_i'')^{(3)}(T_{21},t)$, the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$ :<br><br>$(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$, are positive constants<br>$\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}$ , $\frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}} < 1$ | 115 |
| There exists two constants There exists two constants $(\hat{P}_{20})^{(3)}$ and $(\hat{Q}_{20})^{(3)}$ which together with $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}, (\hat{A}_{20})^{(3)} and\ (\hat{B}_{20})^{(3)}$ and the constants $(a_i)^{(3)}, (a_i')^{(3)}, (b_i)^{(3)}, (b_i')^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}, i = 20,21,22,$ satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(a_i)^{(3)} + (a_i')^{(3)} + (\hat{A}_{20})^{(3)} + (\hat{P}_{20})^{(3)}(\hat{k}_{20})^{(3)}] < 1$<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(b_i)^{(3)} + (b_i')^{(3)} + (\hat{B}_{20})^{(3)} + (\hat{Q}_{20})^{(3)}(\hat{k}_{20})^{(3)}] < 1$ | 116 |
| Where we suppose | |
| $(a_i)^{(4)}, (a_i')^{(4)}, (a_i'')^{(4)}, (b_i)^{(4)}, (b_i')^{(4)}, (b_i'')^{(4)} > 0, \quad i,j = 24,25,26$<br><br>The functions $(a_i'')^{(4)}, (b_i'')^{(4)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(4)}$, $(r_i)^{(4)}$:<br><br>$(a_i'')^{(4)}(T_{25},t) \le (p_i)^{(4)} \le (\hat{A}_{24})^{(4)}$<br><br>$(b_i'')^{(4)}\big((G_{27}),t\big) \le (r_i)^{(4)} \le (b_i')^{(4)} \le (\hat{B}_{24})^{(4)}$ | 117 |
| $lim_{T_2\to\infty}(a_i'')^{(4)}(T_{25},t) = (p_i)^{(4)}$ | 118 |

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|---|---|
| $\lim_{G\to\infty}(b_i'')^{(4)}\left((G_{27}),t\right) = \ (r_i)^{(4)}$<br><br>**Definition of** $(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}$ :<br><br>Where $\boxed{(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}, (p_i)^{(4)}, \ (r_i)^{(4)}}$ are positive constants and $\boxed{i = 24{,}25{,}26}$ | |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i'')^{(4)}(T_{25}',t) - (a_i'')^{(4)}(T_{25},t)\rvert \le (\hat{k}_{24})^{(4)}\lvert T_{25} - T_{25}'\rvert e^{-(\hat{M}_{24})^{(4)}t}$<br><br>$\lvert(b_i'')^{(4)}((G_{27})',t) - (b_i'')^{(4)}\left((G_{27}),t\right)\rvert < (\hat{k}_{24})^{(4)}\lvert\lvert(G_{27}) - (G_{27})'\rvert\rvert e^{-(\hat{M}_{24})^{(4)}t}$ | 119 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(4)}(T_{25}',t)$ and $(a_i'')^{(4)}(T_{25},t)$ . $(T_{25}',t)$ and $(T_{25},t)$ are points belonging to the interval $\left[(\hat{k}_{24})^{(4)}, (\hat{M}_{24})^{(4)}\right]$. It is to be noted that $(a_i'')^{(4)}(T_{25},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{24})^{(4)} = 1$ then the function $(a_i'')^{(4)}(T_{25},t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$ :<br><br>$(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$, are positive constants<br><br>$\frac{(a_i)^{(4)}}{(\hat{M}_{24})^{(4)}}$ , $\frac{(b_i)^{(4)}}{(\hat{M}_{24})^{(4)}} < 1$ | 120 |
| **Definition of** $(\hat{P}_{24})^{(4)}, (\hat{Q}_{24})^{(4)}$ :<br><br>There exists two constants $(\hat{P}_{24})^{(4)}$ and $(\hat{Q}_{24})^{(4)}$ which together with $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}, (\hat{A}_{24})^{(4)} and\ (\hat{B}_{24})^{(4)}$ and the constants $(a_i)^{(4)}, (a_i')^{(4)}, (b_i)^{(4)}, (b_i')^{(4)}, (p_i)^{(4)}, \ (r_i)^{(4)}, i = 24{,}25{,}26$, satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{24})^{(4)}}[(a_i)^{(4)} + (a_i')^{(4)} + \ (\hat{A}_{24})^{(4)} + (\hat{P}_{24})^{(4)}(\hat{k}_{24})^{(4)}] < 1$<br><br>$\frac{1}{(\hat{M}_{24})^{(4)}}[\ (b_i)^{(4)} + (b_i')^{(4)} + \ (\hat{B}_{24})^{(4)} + (\hat{Q}_{24})^{(4)}\ (\hat{k}_{24})^{(4)}] < 1$ | 121 |
| Where we suppose | |
| $(a_i)^{(5)}, (a_i')^{(5)}, (a_i'')^{(5)}, (b_i)^{(5)}, (b_i')^{(5)}, (b_i'')^{(5)} > 0, \quad i,j = 28{,}29{,}30$<br><br>The functions $(a_i'')^{(5)}, (b_i'')^{(5)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(5)}, \ (r_i)^{(5)}$:<br><br>$(a_i'')^{(5)}(T_{29},t) \le (p_i)^{(5)} \le (\hat{A}_{28})^{(5)}$<br><br>$(b_i'')^{(5)}\left((G_{31}),t\right) \le \ (r_i)^{(5)} \le (b_i')^{(5)} \le (\hat{B}_{28})^{(5)}$ | 122 |

$lim_{T_2\to\infty}(a_i'')^{(5)}(T_{29},t) = (p_i)^{(5)}$

$\lim_{G\to\infty}(b_i'')^{(5)}(G_{31},t) = (r_i)^{(5)}$

**Definition of** $(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}$ :

Where $\boxed{(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}, (p_i)^{(5)}, (r_i)^{(5)}}$ are positive constants and $\boxed{i = 28,29,30}$

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They satisfy Lipschitz condition:

$$|(a_i'')^{(5)}(T_{29}',t) - (a_i'')^{(5)}(T_{29},t)| \le (\hat{k}_{28})^{(5)}|T_{29} - T_{29}'|e^{-(\hat{M}_{28})^{(5)}t}$$

$$|(b_i'')^{(5)}((G_{31})',t) - (b_i'')^{(5)}((G_{31}),t)| < (\hat{k}_{28})^{(5)}||(G_{31}) - (G_{31})'||e^{-(\hat{M}_{28})^{(5)}t}$$

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With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(5)}(T_{29}',t)$ and $(a_i'')^{(5)}(T_{29},t)$ . $(T_{29}',t)$ and $(T_{29},t)$ are points belonging to the interval $[(\hat{k}_{28})^{(5)}, (\hat{M}_{28})^{(5)}]$. It is to be noted that $(a_i'')^{(5)}(T_{29},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{28})^{(5)} = 1$ then the function $(a_i'')^{(5)}(T_{29},t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

**Definition of** $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}$ :

$(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)},$ are positive constants

$$\frac{(a_i)^{(5)}}{(\hat{M}_{28})^{(5)}}, \frac{(b_i)^{(5)}}{(\hat{M}_{28})^{(5)}} < 1$$

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**Definition of** $(\hat{P}_{28})^{(5)}, (\hat{Q}_{28})^{(5)}$ :

There exists two constants $(\hat{P}_{28})^{(5)}$ and $(\hat{Q}_{28})^{(5)}$ which together with $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}, (\hat{A}_{28})^{(5)} and\ (\hat{B}_{28})^{(5)}$ and the constants $(a_i)^{(5)}, (a_i')^{(5)}, (b_i)^{(5)}, (b_i')^{(5)}, (p_i)^{(5)}, (r_i)^{(5)}, i = 28,29,30,$ satisfy the inequalities

$$\frac{1}{(\hat{M}_{28})^{(5)}}[(a_i)^{(5)} + (a_i')^{(5)} + (\hat{A}_{28})^{(5)} + (\hat{P}_{28})^{(5)}(\hat{k}_{28})^{(5)}] < 1$$

$$\frac{1}{(\hat{M}_{28})^{(5)}}[(b_i)^{(5)} + (b_i')^{(5)} + (\hat{B}_{28})^{(5)} + (\hat{Q}_{28})^{(5)}(\hat{k}_{28})^{(5)}] < 1$$

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Where we suppose

$(a_i)^{(6)}, (a_i')^{(6)}, (a_i'')^{(6)}, (b_i)^{(6)}, (b_i')^{(6)}, (b_i'')^{(6)} > 0, \quad i,j = 32,33,34$

The functions $(a_i'')^{(6)}, (b_i'')^{(6)}$ are positive continuous increasing and bounded.

**Definition of** $(p_i)^{(6)}, (r_i)^{(6)}$:

$$(a_i'')^{(6)}(T_{33},t) \le (p_i)^{(6)} \le (\hat{A}_{32})^{(6)}$$

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| $(b_i^{''})^{(6)}((G_{35}),t) \leq \ (r_i)^{(6)} \leq (b_i^{'})^{(6)} \leq (\hat{B}_{32})^{(6)}$ | |
| $lim_{T_2\to\infty}(a_i^{''})^{(6)}(T_{33},t) = (p_i)^{(6)}$<br><br>$\lim_{G\to\infty}(b_i^{''})^{(6)}\big((G_{35}),t\big) = \ (r_i)^{(6)}$<br><br>**Definition of** $(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}$ :<br><br>Where $\boxed{(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}, (p_i)^{(6)}, \ (r_i)^{(6)}}$ are positive constants and $\boxed{i = 32,33,34}$ | 128 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(6)}(T_{33}^{'},t) - (a_i^{''})^{(6)}(T_{33},t)\rvert \leq (\hat{k}_{32})^{(6)}\lvert T_{33} - T_{33}^{'}\rvert e^{-(\hat{M}_{32})^{(6)}t}$<br><br>$\lvert(b_i^{''})^{(6)}((G_{35})^{'},t) - (b_i^{''})^{(6)}\big((G_{35}),t\big)\rvert < (\hat{k}_{32})^{(6)}\lvert\lvert(G_{35}) - (G_{35})^{'}\rvert\rvert e^{-(\hat{M}_{32})^{(6)}t}$ | |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(6)}(T_{33}^{'},t)$ and$(a_i^{''})^{(6)}(T_{33},t)$ . $(T_{33}^{'},t)$ and $(T_{33},t)$ are points belonging to the interval $\big[(\hat{k}_{32})^{(6)}, (\hat{M}_{32})^{(6)}\big]$. It is to be noted that $(a_i^{''})^{(6)}(T_{33},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{32})^{(6)} = 1$ then the function $(a_i^{''})^{(6)}(T_{33},t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}$ :<br><br>$(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}$, are positive constants<br><br>$\frac{(a_i)^{(6)}}{(\hat{M}_{32})^{(6)}}$ , $\frac{(b_i)^{(6)}}{(\hat{M}_{32})^{(6)}} < 1$ | 129 |
| **Definition of** $(\hat{P}_{32})^{(6)}, (\hat{Q}_{32})^{(6)}$ :<br><br>There exists two constants $(\hat{P}_{32})^{(6)}$ and $(\hat{Q}_{32})^{(6)}$ which together with $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}, (\hat{A}_{32})^{(6)} and\ (\hat{B}_{32})^{(6)}$ and the constants $(a_i)^{(6)}, (a_i^{'})^{(6)}, (b_i)^{(6)}, (b_i^{'})^{(6)}, (p_i)^{(6)}, \ (r_i)^{(6)}, i = 32,33,34,$ satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{32})^{(6)}}[\,(a_i)^{(6)} + (a_i^{'})^{(6)} + \ (\hat{A}_{32})^{(6)} + \ (\hat{P}_{32})^{(6)}\,(\hat{k}_{32})^{(6)}] < 1$<br><br>$\frac{1}{(\hat{M}_{32})^{(6)}}[\ (b_i)^{(6)} + (b_i^{'})^{(6)} + \ (\hat{B}_{32})^{(6)} + \ (\hat{Q}_{32})^{(6)}\ (\hat{k}_{32})^{(6)}] < 1$ | 130 |
| Where we suppose | |
| (A) $(a_i)^{(7)}, (a_i^{'})^{(7)}, (a_i^{''})^{(7)}, (b_i)^{(7)}, (b_i^{'})^{(7)}, (b_i^{''})^{(7)} > 0, \quad i,j = 36,37,38$<br><br>(B) The functions $(a_i^{''})^{(7)}, (b_i^{''})^{(7)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(7)}, \ (r_i)^{(7)}$: | 131 |

$(a_i^{''})^{(7)}(T_{37},t) \le (p_i)^{(7)} \le (\hat{A}_{36})^{(7)}$

$(b_i^{''})^{(7)}(G_{39},t) \le (r_i)^{(7)} \le (b_i^{'})^{(7)} \le (\hat{B}_{36})^{(7)}$

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*(C)* $\lim_{T_2\to\infty} (a_i^{''})^{(7)}(T_{37},t) = (p_i)^{(7)}$

*(D)*

$\lim_{G\to\infty}(b_i^{''})^{(7)}\left((G_{39}),t\right) = (r_i)^{(7)}$

**Definition of** $(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}$ :

Where $\boxed{(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}, (p_i)^{(7)}, (r_i)^{(7)}}$ are positive constants and $\boxed{i = 36,37,38}$

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They satisfy Lipschitz condition:

$|(a_i^{''})^{(7)}(T_{37}^{'},t) - (a_i^{''})^{(7)}(T_{37},t)| \le (\hat{k}_{36})^{(7)}|T_{37} - T_{37}^{'}|e^{-(\hat{M}_{36})^{(7)}t}$

$|(b_i^{''})^{(7)}((G_{39})^{'},t) - (b_i^{''})^{(7)}\left((G_{39}),t\right)| < (\hat{k}_{36})^{(7)}||(G_{39}) - (G_{39})^{'}||e^{-(\hat{M}_{36})^{(7)}t}$

With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(7)}(T_{37}^{'},t)$ and$(a_i^{''})^{(7)}(T_{37},t)$ . $(T_{37}^{'},t)$ and $(T_{37},t)$ are points belonging to the interval $\left[(\hat{k}_{36})^{(7)}, (\hat{M}_{36})^{(7)}\right]$. It is to be noted that $(a_i^{''})^{(7)}(T_{37},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{36})^{(7)} = 1$ then the function $(a_i^{''})^{(7)}(T_{37},t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

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**Definition of** $(\hat{M}_{36})^{(7)}, (\hat{k}_{36})^{(7)}$ :

(E) $(\hat{M}_{36})^{(7)}, (\hat{k}_{36})^{(7)}$, are positive constants

$\frac{(a_i)^{(7)}}{(\hat{M}_{36})^{(7)}}, \frac{(b_i)^{(7)}}{(\hat{M}_{36})^{(7)}} < 1$

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**Definition of** $(\hat{P}_{36})^{(7)}, (\hat{Q}_{36})^{(7)}$ :

(F) There exists two constants $(\hat{P}_{36})^{(7)}$ and $(\hat{Q}_{36})^{(7)}$ which together with $(\hat{M}_{36})^{(7)}, (\hat{k}_{36})^{(7)}, (\hat{A}_{36})^{(7)} and\ (\hat{B}_{36})^{(7)}$ and the constants $(a_i)^{(7)}, (a_i^{'})^{(7)}, (b_i)^{(7)}, (b_i^{'})^{(7)}, (p_i)^{(7)}, (r_i)^{(7)}, i = 36,37,38$, satisfy the inequalities

$\frac{1}{(\hat{M}_{36})^{(7)}}[(a_i)^{(7)} + (a_i^{'})^{(7)} + (\hat{A}_{36})^{(7)} + (\hat{P}_{36})^{(7)}(\hat{k}_{36})^{(7)}] < 1$

$\frac{1}{(\hat{M}_{36})^{(7)}}[(b_i)^{(7)} + (b_i^{'})^{(7)} + (\hat{B}_{36})^{(7)} + (\hat{Q}_{36})^{(7)}(\hat{k}_{36})^{(7)}] < 1$

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| Where we suppose | |
| $(a_i)^{(8)}, (a_i^{'})^{(8)}, (a_i^{''})^{(8)}, (b_i)^{(8)}, (b_i^{'})^{(8)}, (b_i^{''})^{(8)} > 0, \quad i,j = 40,41,42$ | 136 |
| The functions $(a_i^{''})^{(8)}, (b_i^{''})^{(8)}$ are positive continuous increasing and bounded | |
| **Definition of** $(p_i)^{(8)}$, $(r_i)^{(8)}$: | 137 |
| $(a_i^{''})^{(8)}(T_{41}, t) \leq (p_i)^{(8)} \leq (\hat{A}_{40})^{(8)}$ | 138 |
| $(b_i^{''})^{(8)}((G_{43}), t) \leq (r_i)^{(8)} \leq (b_i^{'})^{(8)} \leq (\hat{B}_{40})^{(8)}$ | 139 |
| $\lim_{T_2 \to \infty} (a_i^{''})^{(8)}(T_{41}, t) = (p_i)^{(8)}$ | 140 |
| $\lim_{G \to \infty} (b_i^{''})^{(8)}((G_{43}), t) = (r_i)^{(8)}$ | 141 |
| **Definition of** $(\hat{A}_{40})^{(8)}, (\hat{B}_{40})^{(8)}$ :<br>Where $\boxed{(\hat{A}_{40})^{(8)}, (\hat{B}_{40})^{(8)}, (p_i)^{(8)}, (r_i)^{(8)}}$ are positive constants and $\boxed{i = 40,41,42}$ | |
| They satisfy Lipschitz condition: | |
| $\lvert(a_i^{''})^{(8)}(T_{41}^{'}, t) - (a_i^{''})^{(8)}(T_{41}, t)\rvert \leq (\hat{k}_{40})^{(8)} \lvert T_{41} - T_{41}^{'}\rvert e^{-(\hat{M}_{40})^{(8)} t}$ | 142 |
| $\lvert(b_i^{''})^{(8)}((G_{43})^{'}, t) - (b_i^{''})^{(8)}((G_{43}), t)\rvert < (\hat{k}_{40})^{(8)} \lvert\lvert(G_{43}) - (G_{43})^{'}\rvert\rvert e^{-(\hat{M}_{40})^{(8)} t}$ | 143 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(8)}(T_{41}^{'}, t)$ and $(a_i^{''})^{(8)}(T_{41}, t)$ . $(T_{41}^{'}, t)$ and $(T_{41}, t)$ are points belonging to the interval $\left[(\hat{k}_{40})^{(8)}, (\hat{M}_{40})^{(8)}\right]$. It is to be noted that $(a_i^{''})^{(8)}(T_{41}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{40})^{(8)} = 1$ then the function $(a_i^{''})^{(8)}(T_{41}, t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}$ : | |
| $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}$, are positive constants | |
| $\frac{(a_i)^{(8)}}{(\hat{M}_{40})^{(8)}}$ , $\frac{(b_i)^{(8)}}{(\hat{M}_{40})^{(8)}} < 1$ | 144 |
| **Definition of** $(\hat{P}_{40})^{(8)}, (\hat{Q}_{40})^{(8)}$ :<br>There exists two constants $(\hat{P}_{40})^{(8)}$ and $(\hat{Q}_{40})^{(8)}$ which together with $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}, (\hat{A}_{40})^{(8)}$ $(\hat{B}_{40})^{(8)}$ and the constants $(a_i)^{(8)}, (a_i^{'})^{(8)}, (b_i)^{(8)}, (b_i^{'})^{(8)}, (p_i)^{(8)}$, $(r_i)^{(8)}, i = 40,41,42,$<br>Satisfy the inequalities | |

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| $\frac{1}{(\hat{M}_{40})^{(8)}}[(a_i)^{(8)} + (a_i')^{(8)} + (\hat{A}_{40})^{(8)} + (\hat{P}_{40})^{(8)}(\hat{k}_{40})^{(8)}] < 1$ | 145 |
| $\frac{1}{(\hat{M}_{40})^{(8)}}[(b_i)^{(8)} + (b_i')^{(8)} + (\hat{B}_{40})^{(8)} + (\hat{Q}_{40})^{(8)}(\hat{k}_{40})^{(8)}] < 1$ | 146 |
| Where we suppose | |
| $(a_i)^{(9)}, (a_i')^{(9)}, (a_i'')^{(9)}, (b_i)^{(9)}, (b_i')^{(9)}, (b_i'')^{(9)} > 0, \quad i,j = 44,45,46$<br>The functions $(a_i'')^{(9)}, (b_i'')^{(9)}$ are positive continuous increasing and bounded.<br>**Definition of** $(p_i)^{(9)}$, $(r_i)^{(9)}$:<br>$(a_i'')^{(9)}(T_{45},t) \le (p_i)^{(9)} \le (\hat{A}_{44})^{(9)}$<br>$(b_i'')^{(9)}(G_{47},t) \le (r_i)^{(9)} \le (b_i')^{(9)} \le (\hat{B}_{44})^{(9)}$ | 146 A |
| $\lim_{T_2\to\infty}(a_i'')^{(9)}(T_{45},t) = (p_i)^{(9)}$<br>$\lim_{G\to\infty}(b_i'')^{(9)}(G_{47},t) = (r_i)^{(9)}$<br>**Definition of** $(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}$ :<br>Where $\boxed{(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}, (p_i)^{(9)}, (r_i)^{(9)}}$ are positive constants and $\boxed{i = 44,45,46}$ | |
| They satisfy Lipschitz condition:<br>$\lvert(a_i'')^{(9)}(T_{45}',t) - (a_i'')^{(9)}(T_{45},t)\rvert \le (\hat{k}_{44})^{(9)}\lvert T_{45} - T_{45}'\rvert e^{-(\hat{M}_{44})^{(9)}t}$<br>$\lvert(b_i'')^{(9)}((G_{47})',t) - (b_i'')^{(9)}((G_{47}),t)\rvert < (\hat{k}_{44})^{(9)}\lvert\lvert(G_{47}) - (G_{47})'\rvert\rvert e^{-(\hat{M}_{44})^{(9)}t}$ | |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(9)}(T_{45}',t)$ and $(a_i'')^{(9)}(T_{45},t)$ . $(T_{45}',t)$ and $(T_{45},t)$ are points belonging to the interval $[(\hat{k}_{44})^{(9)}, (\hat{M}_{44})^{(9)}]$. It is to be noted that $(a_i'')^{(9)}(T_{45},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{44})^{(9)} = 1$ then the function $(a_i'')^{(9)}(T_{45},t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}$ :<br>$(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}$, are positive constants<br>$\frac{(a_i)^{(9)}}{(\hat{M}_{44})^{(9)}}, \frac{(b_i)^{(9)}}{(\hat{M}_{44})^{(9)}} < 1$ | |
| **Definition of** $(\hat{P}_{44})^{(9)}, (\hat{Q}_{44})^{(9)}$ :<br>There exists two constants $(\hat{P}_{44})^{(9)}$ and $(\hat{Q}_{44})^{(9)}$ which together with $(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}, (\hat{A}_{44})^{(9)} and (\hat{B}_{44})^{(9)}$ and the constants | |

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| $(a_i)^{(9)}, (a_i')^{(9)}, (b_i)^{(9)}, (b_i')^{(9)}, (p_i)^{(9)}, \ (r_i)^{(9)}, i = 44,45,46,$ satisfy the inequalities<br>$\frac{1}{(\hat{M}_{44})^{(9)}}[(a_i)^{(9)} + (a_i')^{(9)} + (\hat{A}_{44})^{(9)} + (\hat{P}_{44})^{(9)} (\hat{k}_{44})^{(9)}] < 1$<br>$\frac{1}{(\hat{M}_{44})^{(9)}}[(b_i)^{(9)} + (b_i')^{(9)} + (\hat{B}_{44})^{(9)} + (\hat{Q}_{44})^{(9)} (\hat{k}_{44})^{(9)}] < 1$ | |
| **Theorem 1:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>Definition of $G_i(0), T_i(0)$ :<br>$G_i(t) \le (\hat{P}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le (\hat{Q}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 147 |
| **Theorem 2 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$<br>$G_i(t) \le (\hat{P}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ , $G_i(0) = G_i^0 > 0$<br>$T_i(t) \le (\hat{Q}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ , $T_i(0) = T_i^0 > 0$ | 148 |
| **Theorem 3 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>$G_i(t) \le (\hat{P}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ , $G_i(0) = G_i^0 > 0$<br>$T_i(t) \le (\hat{Q}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ , $T_i(0) = T_i^0 > 0$ | 149 |
| **Theorem 4 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \le (\hat{P}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le (\hat{Q}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 150 |
| **Theorem 5 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \le (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le (\hat{Q}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 151 |

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| **Theorem 6 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br><br>**Definition of** $G_i(0)\,, T_i(0)$ :<br><br>$G_i(t) \le \left(\hat{P}_{32}\right)^{(6)} e^{(\hat{M}_{32})^{(6)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br><br>$T_i(t) \le \left(\hat{Q}_{32}\right)^{(6)} e^{(\hat{M}_{32})^{(6)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 152 |
| **Theorem 7:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br><br>**Definition of** $G_i(0)\,, T_i(0)$ :<br><br>$G_i(t) \le \left(\hat{P}_{36}\right)^{(7)} e^{(\hat{M}_{36})^{(7)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{36}\right)^{(7)} e^{(\hat{M}_{36})^{(7)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 |
| **Theorem 8:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br><br>**Definition of** $G_i(0)\,, T_i(0)$ :<br><br>$G_i(t) \le \left(\hat{P}_{40}\right)^{(8)} e^{(\hat{M}_{40})^{(8)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{40}\right)^{(8)} e^{(\hat{M}_{40})^{(8)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 A |
| **Theorem 9:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br><br>**Definition of** $G_i(0)\,, T_i(0)$ :<br><br>$G_i(t) \le \left(\hat{P}_{44}\right)^{(9)} e^{(\hat{M}_{44})^{(9)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br><br>$T_i(t) \le \left(\hat{Q}_{44}\right)^{(9)} e^{(\hat{M}_{44})^{(9)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 B |
| **Proof:** Consider operator $\mathcal{A}^{(1)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i$: $\mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 154 |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \le \left(\hat{P}_{13}\right)^{(1)}$ , $T_i^0 \le \left(\hat{Q}_{13}\right)^{(1)}$, | 155 |
| $0 \le G_i(t) - G_i^0 \le \left(\hat{P}_{13}\right)^{(1)} e^{(\hat{M}_{13})^{(1)}t}$ | 156 |
| $0 \le T_i(t) - T_i^0 \le \left(\hat{Q}_{13}\right)^{(1)} e^{(\hat{M}_{13})^{(1)}t}$ | 157 |
| By<br><br>$\bar{G}_{13}(t) = G_{13}^0 + \int_0^t \left[(a_{13})^{(1)} G_{14}\left(s_{(13)}\right) - \left((a_{13}')^{(1)} + a_{13}'')^{(1)}\left(T_{14}\left(s_{(13)}\right), s_{(13)}\right)\right) G_{13}\left(s_{(13)}\right)\right] ds_{(13)}$ | 158 |

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| $\bar{G}_{14}(t) = G_{14}^0 + \int_0^t \left[(a_{14})^{(1)} G_{13}(s_{(13)}) - \left((a_{14}')^{(1)} + (a_{14}'')^{(1)}(T_{14}(s_{(13)}), s_{(13)})\right) G_{14}(s_{(13)})\right] ds_{(13)}$ | |
| $\bar{G}_{15}(t) = G_{15}^0 + \int_0^t \left[(a_{15})^{(1)} G_{14}(s_{(13)}) - \left((a_{15}')^{(1)} + (a_{15}'')^{(1)}(T_{14}(s_{(13)}), s_{(13)})\right) G_{15}(s_{(13)})\right] ds_{(13)}$ | |
| $\bar{T}_{13}(t) = T_{13}^0 + \int_0^t \left[(b_{13})^{(1)} T_{14}(s_{(13)}) - \left((b_{13}')^{(1)} - (b_{13}'')^{(1)}(G(s_{(13)}), s_{(13)})\right) T_{13}(s_{(13)})\right] ds_{(13)}$ | |
| $\bar{T}_{14}(t) = T_{14}^0 + \int_0^t \left[(b_{14})^{(1)} T_{13}(s_{(13)}) - \left((b_{14}')^{(1)} - (b_{14}'')^{(1)}(G(s_{(13)}), s_{(13)})\right) T_{14}(s_{(13)})\right] ds_{(13)}$ | |
| $\bar{T}_{15}(t) = T_{15}^0 + \int_0^t \left[(b_{15})^{(1)} T_{14}(s_{(13)}) - \left((b_{15}')^{(1)} - (b_{15}'')^{(1)}(G(s_{(13)}), s_{(13)})\right) T_{15}(s_{(13)})\right] ds_{(13)}$<br>Where $s_{(13)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(2)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 159 |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{16})^{(2)}$, $T_i^0 \le (\hat{Q}_{16})^{(2)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ | |
| By<br>$\bar{G}_{16}(t) = G_{16}^0 + \int_0^t \left[(a_{16})^{(2)} G_{17}(s_{(16)}) - \left((a_{16}')^{(2)} + a_{16}'')^{(2)}(T_{17}(s_{(16)}), s_{(16)})\right) G_{16}(s_{(16)})\right] ds_{(16)}$ | 160 |
| $\bar{G}_{17}(t) = G_{17}^0 + \int_0^t \left[(a_{17})^{(2)} G_{16}(s_{(16)}) - \left((a_{17}')^{(2)} + (a_{17}'')^{(2)}(T_{17}(s_{(16)}), s_{(17)})\right) G_{17}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{G}_{18}(t) = G_{18}^0 + \int_0^t \left[(a_{18})^{(2)} G_{17}(s_{(16)}) - \left((a_{18}')^{(2)} + (a_{18}'')^{(2)}(T_{17}(s_{(16)}), s_{(16)})\right) G_{18}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{T}_{16}(t) = T_{16}^0 + \int_0^t \left[(b_{16})^{(2)} T_{17}(s_{(16)}) - \left((b_{16}')^{(2)} - (b_{16}'')^{(2)}(G_{19}(s_{(16)}), s_{(16)})\right) T_{16}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{T}_{17}(t) = T_{17}^0 + \int_0^t \left[(b_{17})^{(2)} T_{16}(s_{(16)}) - \left((b_{17}')^{(2)} - (b_{17}'')^{(2)}(G_{19}(s_{(16)}), s_{(16)})\right) T_{17}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{T}_{18}(t) = T_{18}^0 + \int_0^t \left[(b_{18})^{(2)} T_{17}(s_{(16)}) - \left((b_{18}')^{(2)} - (b_{18}'')^{(2)}(G_{19}(s_{(16)}), s_{(16)})\right) T_{18}(s_{(16)})\right] ds_{(16)}$<br>Where $s_{(16)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(3)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |

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| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{20})^{(3)}$ , $T_i^0 \leq (\hat{Q}_{20})^{(3)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{20})^{(3)} e^{(\bar{M}_{20})^{(3)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{20})^{(3)} e^{(\bar{M}_{20})^{(3)} t}$ | |
| By<br>$\bar{G}_{20}(t) = G_{20}^0 + \int_0^t \left[ (a_{20})^{(3)} G_{21}(s_{(20)}) - \left( (a_{20}^{'})^{(3)} + a_{20}^{''})^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{20}(s_{(20)}) \right] ds_{(20)}$ | 161 |
| $\bar{G}_{21}(t) = G_{21}^0 + \int_0^t \left[ (a_{21})^{(3)} G_{20}(s_{(20)}) - \left( (a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{21}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{G}_{22}(t) = G_{22}^0 + \int_0^t \left[ (a_{22})^{(3)} G_{21}(s_{(20)}) - \left( (a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{22}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{20}(t) = T_{20}^0 + \int_0^t \left[ (b_{20})^{(3)} T_{21}(s_{(20)}) - \left( (b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{20}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{21}(t) = T_{21}^0 + \int_0^t \left[ (b_{21})^{(3)} T_{20}(s_{(20)}) - \left( (b_{21}^{'})^{(3)} - (b_{21}^{''})^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{21}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{22}(t) = T_{22}^0 + \int_0^t \left[ (b_{22})^{(3)} T_{21}(s_{(20)}) - \left( (b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{22}(s_{(20)}) \right] ds_{(20)}$<br>Where $s_{(20)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:** Consider operator $\mathcal{A}^{(4)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i$: $\mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{24})^{(4)}$ , $T_i^0 \leq (\hat{Q}_{24})^{(4)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{24})^{(4)} e^{(\bar{M}_{24})^{(4)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{24})^{(4)} e^{(\bar{M}_{24})^{(4)} t}$ | |
| By<br>$\bar{G}_{24}(t) = G_{24}^0 + \int_0^t \left[ (a_{24})^{(4)} G_{25}(s_{(24)}) - \left( (a_{24}^{'})^{(4)} + a_{24}^{''})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{24}(s_{(24)}) \right] ds_{(24)}$ | 162 |
| $\bar{G}_{25}(t) = G_{25}^0 + \int_0^t \left[ (a_{25})^{(4)} G_{24}(s_{(24)}) - \left( (a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{25}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{G}_{26}(t) = G_{26}^0 + \int_0^t \left[ (a_{26})^{(4)} G_{25}(s_{(24)}) - \left( (a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{26}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{24}(t) = T_{24}^0 + \int_0^t \left[ (b_{24})^{(4)} T_{25}(s_{(24)}) - \left( (b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{24}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{25}(t) = T_{25}^0 + \int_0^t \left[ (b_{25})^{(4)} T_{24}(s_{(24)}) - \left( (b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{25}(s_{(24)}) \right] ds_{(24)}$ | |

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| $\overline{T}_{26}(t) = T_{26}^0 + \int_0^t \left[(b_{26})^{(4)} T_{25}(s_{(24)}) - \left((b_{26}')^{(4)} - (b_{26}'')^{(4)}(G_{27}(s_{(24)}), s_{(24)})\right) T_{26}(s_{(24)})\right] ds_{(24)}$<br>Where $s_{(24)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:** Consider operator $\mathcal{A}^{(5)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \le (\hat{P}_{28})^{(5)}$ , $T_i^0 \le (\hat{Q}_{28})^{(5)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| By<br>$\bar{G}_{28}(t) = G_{28}^0 + \int_0^t \left[(a_{28})^{(5)} G_{29}(s_{(28)}) - \left((a_{28}')^{(5)} + a_{28}'')^{(5)}(T_{29}(s_{(28)}), s_{(28)})\right) G_{28}(s_{(28)})\right] ds_{(28)}$ | 163 |
| $\bar{G}_{29}(t) = G_{29}^0 + \int_0^t \left[(a_{29})^{(5)} G_{28}(s_{(28)}) - \left((a_{29}')^{(5)} + (a_{29}'')^{(5)}(T_{29}(s_{(28)}), s_{(28)})\right) G_{29}(s_{(28)})\right] ds_{(28)}$ | |
| $\bar{G}_{30}(t) = G_{30}^0 + \int_0^t \left[(a_{30})^{(5)} G_{29}(s_{(28)}) - \left((a_{30}')^{(5)} + (a_{30}'')^{(5)}(T_{29}(s_{(28)}), s_{(28)})\right) G_{30}(s_{(28)})\right] ds_{(28)}$ | |
| $\overline{T}_{28}(t) = T_{28}^0 + \int_0^t \left[(b_{28})^{(5)} T_{29}(s_{(28)}) - \left((b_{28}')^{(5)} - (b_{28}'')^{(5)}(G_{31}(s_{(28)}), s_{(28)})\right) T_{28}(s_{(28)})\right] ds_{(28)}$ | |
| $\overline{T}_{29}(t) = T_{29}^0 + \int_0^t \left[(b_{29})^{(5)} T_{28}(s_{(28)}) - \left((b_{29}')^{(5)} - (b_{29}'')^{(5)}(G_{31}(s_{(28)}), s_{(28)})\right) T_{29}(s_{(28)})\right] ds_{(28)}$ | |
| $\overline{T}_{30}(t) = T_{30}^0 + \int_0^t \left[(b_{30})^{(5)} T_{29}(s_{(28)}) - \left((b_{30}')^{(5)} - (b_{30}'')^{(5)}(G_{31}(s_{(28)}), s_{(28)})\right) T_{30}(s_{(28)})\right] ds_{(28)}$<br>Where $s_{(28)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(6)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \le (\hat{P}_{32})^{(6)}$ , $T_i^0 \le (\hat{Q}_{32})^{(6)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ | |
| By<br>$\bar{G}_{32}(t) = G_{32}^0 + \int_0^t \left[(a_{32})^{(6)} G_{33}(s_{(32)}) - \left((a_{32}')^{(6)} + a_{32}'')^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{32}(s_{(32)})\right] ds_{(32)}$ | 164 |
| $\bar{G}_{33}(t) = G_{33}^0 + \int_0^t \left[(a_{33})^{(6)} G_{32}(s_{(32)}) - \left((a_{33}')^{(6)} + (a_{33}'')^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{33}(s_{(32)})\right] ds_{(32)}$ | |

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| $\bar{G}_{34}(t) = G_{34}^0 + \int_0^t \left[(a_{34})^{(6)} G_{33}(s_{(32)}) - \left((a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{34}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{T}_{32}(t) = T_{32}^0 + \int_0^t \left[(b_{32})^{(6)} T_{33}(s_{(32)}) - \left((b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{32}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{T}_{33}(t) = T_{33}^0 + \int_0^t \left[(b_{33})^{(6)} T_{32}(s_{(32)}) - \left((b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{33}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{T}_{34}(t) = T_{34}^0 + \int_0^t \left[(b_{34})^{(6)} T_{33}(s_{(32)}) - \left((b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{34}(s_{(32)})\right] ds_{(32)}$<br>Where $s_{(32)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(7)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{36})^{(7)}$, $T_i^0 \le (\hat{Q}_{36})^{(7)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ | |
| By<br>$\bar{G}_{36}(t) = G_{36}^0 + \int_0^t \left[(a_{36})^{(7)} G_{37}(s_{(36)}) - \left((a_{36}^{'})^{(7)} + a_{36}^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{36}(s_{(36)})\right] ds_{(36)}$ | 165 |
| $\bar{G}_{37}(t) = G_{37}^0 + \int_0^t \left[(a_{37})^{(7)} G_{36}(s_{(36)}) - \left((a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{37}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{G}_{38}(t) = G_{38}^0 + \int_0^t \left[(a_{38})^{(7)} G_{37}(s_{(36)}) - \left((a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{38}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{36}(t) = T_{36}^0 + \int_0^t \left[(b_{36})^{(7)} T_{37}(s_{(36)}) - \left((b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{36}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{37}(t) = T_{37}^0 + \int_0^t \left[(b_{37})^{(7)} T_{36}(s_{(36)}) - \left((b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{37}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{38}(t) = T_{38}^0 + \int_0^t \left[(b_{38})^{(7)} T_{37}(s_{(36)}) - \left((b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{38}(s_{(36)})\right] ds_{(36)}$<br>Where $s_{(36)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(8)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ | |

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| which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{40})^{(8)}$ , $T_i^0 \leq (\hat{Q}_{40})^{(8)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{40})^{(8)} e^{(\bar{M}_{40})^{(8)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{40})^{(8)} e^{(\bar{M}_{40})^{(8)} t}$ | |
| By<br>$\bar{G}_{40}(t) = G_{40}^0 + \int_0^t \left[ (a_{40})^{(8)} G_{41}(s_{(40)}) - \left( (a'_{40})^{(8)} + a''_{40})^{(8)} (T_{41}(s_{(40)}), s_{(40)}) \right) G_{40}(s_{(40)}) \right] ds_{(40)}$ | 166 |
| $\bar{G}_{41}(t) = G_{41}^0 + \int_0^t \left[ (a_{41})^{(8)} G_{40}(s_{(40)}) - \left( (a'_{41})^{(8)} + (a''_{41})^{(8)} (T_{41}(s_{(40)}), s_{(40)}) \right) G_{41}(s_{(40)}) \right] ds_{(40)}$ | |
| $\bar{G}_{42}(t) = G_{42}^0 + \int_0^t \left[ (a_{42})^{(8)} G_{41}(s_{(40)}) - \left( (a'_{42})^{(8)} + (a''_{42})^{(8)} (T_{41}(s_{(40)}), s_{(40)}) \right) G_{42}(s_{(40)}) \right] ds_{(40)}$ | |
| $\bar{T}_{40}(t) = T_{40}^0 + \int_0^t \left[ (b_{40})^{(8)} T_{41}(s_{(40)}) - \left( (b'_{40})^{(8)} - (b''_{40})^{(8)} (G_{43}(s_{(40)}), s_{(40)}) \right) T_{40}(s_{(40)}) \right] ds_{(40)}$ | |
| $\bar{T}_{41}(t) = T_{41}^0 + \int_0^t \left[ (b_{41})^{(8)} T_{40}(s_{(40)}) - \left( (b'_{41})^{(8)} - (b''_{41})^{(8)} (G_{43}(s_{(40)}), s_{(40)}) \right) T_{41}(s_{(40)}) \right] ds_{(40)}$ | |
| $\bar{T}_{42}(t) = T_{42}^0 + \int_0^t \left[ (b_{42})^{(8)} T_{41}(s_{(40)}) - \left( (b'_{42})^{(8)} - (b''_{42})^{(8)} (G_{43}(s_{(40)}), s_{(40)}) \right) T_{42}(s_{(40)}) \right] ds_{(40)}$<br>Where $s_{(40)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(9)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 166 A |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{44})^{(9)}$ , $T_i^0 \leq (\hat{Q}_{44})^{(9)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{44})^{(9)} e^{(\bar{M}_{44})^{(9)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{44})^{(9)} e^{(\bar{M}_{44})^{(9)} t}$ | |
| By<br>$\bar{G}_{44}(t) = G_{44}^0 + \int_0^t \left[ (a_{44})^{(9)} G_{45}(s_{(44)}) - \left( (a'_{44})^{(9)} + a''_{44})^{(9)} (T_{45}(s_{(44)}), s_{(44)}) \right) G_{44}(s_{(44)}) \right] ds_{(44)}$ | |
| $\bar{G}_{45}(t) = G_{45}^0 + \int_0^t \left[ (a_{45})^{(9)} G_{44}(s_{(44)}) - \left( (a'_{45})^{(9)} + (a''_{45})^{(9)} (T_{45}(s_{(44)}), s_{(44)}) \right) G_{45}(s_{(44)}) \right] ds_{(44)}$ | |
| $\bar{G}_{46}(t) = G_{46}^0 + \int_0^t \left[ (a_{46})^{(9)} G_{45}(s_{(44)}) - \left( (a'_{46})^{(9)} + (a''_{46})^{(9)} (T_{45}(s_{(44)}), s_{(44)}) \right) G_{46}(s_{(44)}) \right] ds_{(44)}$ | |

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| $\bar{T}_{44}(t) = T_{44}^0 + \int_0^t \left[(b_{44})^{(9)} T_{45}(s_{(44)}) - \left((b'_{44})^{(9)} - (b''_{44})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{44}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{45}(t) = T_{45}^0 + \int_0^t \left[(b_{45})^{(9)} T_{44}(s_{(44)}) - \left((b'_{45})^{(9)} - (b''_{45})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{45}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{46}(t) = T_{46}^0 + \int_0^t \left[(b_{46})^{(9)} T_{45}(s_{(44)}) - \left((b'_{46})^{(9)} - (b''_{46})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{46}(s_{(44)})\right] ds_{(44)}$<br>Where $s_{(44)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| The operator $\mathcal{A}^{(1)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{13}(t) \le G_{13}^0 + \int_0^t \left[(a_{13})^{(1)} \left(G_{14}^0 + (\hat{P}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} s_{(13)}}\right)\right] ds_{(13)} =$<br>$\left(1 + (a_{13})^{(1)} t\right) G_{14}^0 + \frac{(a_{13})^{(1)} (\hat{P}_{13})^{(1)}}{(\hat{M}_{13})^{(1)}} \left(e^{(\hat{M}_{13})^{(1)} t} - 1\right)$ | 167 |
| From which it follows that<br>$(G_{13}(t) - G_{13}^0) e^{-(\hat{M}_{13})^{(1)} t} \le \frac{(a_{13})^{(1)}}{(\hat{M}_{13})^{(1)}} \left[\left((\hat{P}_{13})^{(1)} + G_{14}^0\right) e^{\left(-\frac{(\hat{P}_{13})^{(1)} + G_{14}^0}{G_{14}^0}\right)} + (\hat{P}_{13})^{(1)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 1 | 168 |
| Analogous inequalities hold also for $G_{14}, G_{15}, T_{13}, T_{14}, T_{15}$ | |
| The operator $\mathcal{A}^{(2)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{16}(t) \le G_{16}^0 + \int_0^t \left[(a_{16})^{(2)} \left(G_{17}^0 + (\hat{P}_{16})^{(6)} e^{(\hat{M}_{16})^{(2)} s_{(16)}}\right)\right] ds_{(16)} =$<br>$\left(1 + (a_{16})^{(2)} t\right) G_{17}^0 + \frac{(a_{16})^{(2)} (\hat{P}_{16})^{(2)}}{(\hat{M}_{16})^{(2)}} \left(e^{(\hat{M}_{16})^{(2)} t} - 1\right)$ | 169 |
| From which it follows that<br>$(G_{16}(t) - G_{16}^0) e^{-(\hat{M}_{16})^{(2)} t} \le \frac{(a_{16})^{(2)}}{(\hat{M}_{16})^{(2)}} \left[\left((\hat{P}_{16})^{(2)} + G_{17}^0\right) e^{\left(-\frac{(\hat{P}_{16})^{(2)} + G_{17}^0}{G_{17}^0}\right)} + (\hat{P}_{16})^{(2)}\right]$ | 170 |
| Analogous inequalities hold also for $G_{17}, G_{18}, T_{16}, T_{17}, T_{18}$ | |
| The operator $\mathcal{A}^{(3)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{20}(t) \le G_{20}^0 + \int_0^t \left[(a_{20})^{(3)} \left(G_{21}^0 + (\hat{P}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} s_{(20)}}\right)\right] ds_{(20)} =$<br>$\left(1 + (a_{20})^{(3)} t\right) G_{21}^0 + \frac{(a_{20})^{(3)} (\hat{P}_{20})^{(3)}}{(\hat{M}_{20})^{(3)}} \left(e^{(\hat{M}_{20})^{(3)} t} - 1\right)$ | 171 |
| From which it follows that | 172 |

$$(G_{20}(t)-G_{20}^0)e^{-(\hat{M}_{20})^{(3)}t} \le \frac{(a_{20})^{(3)}}{(\hat{M}_{20})^{(3)}}\left[((\hat{P}_{20})^{(3)}+G_{21}^0)e^{\left(-\frac{(\hat{P}_{20})^{(3)}+G_{21}^0}{G_{21}^0}\right)}+(\hat{P}_{20})^{(3)}\right]$$

Analogous inequalities hold also for $G_{21}, G_{22}, T_{20}, T_{21}, T_{22}$

The operator $\mathcal{A}^{(4)}$ maps the space of functions satisfying into itself .Indeed it is obvious that 173

$$G_{24}(t) \le G_{24}^0 + \int_0^t \left[(a_{24})^{(4)}\left(G_{25}^0+(\hat{P}_{24})^{(4)}e^{(\hat{M}_{24})^{(4)}s_{(24)}}\right)\right] ds_{(24)} =$$

$$\left(1+(a_{24})^{(4)}t\right)G_{25}^0 + \frac{(a_{24})^{(4)}(\hat{P}_{24})^{(4)}}{(\hat{M}_{24})^{(4)}}\left(e^{(\hat{M}_{24})^{(4)}t}-1\right)$$

From which it follows that 174

$$(G_{24}(t)-G_{24}^0)e^{-(\hat{M}_{24})^{(4)}t} \le \frac{(a_{24})^{(4)}}{(\hat{M}_{24})^{(4)}}\left[((\hat{P}_{24})^{(4)}+G_{25}^0)e^{\left(-\frac{(\hat{P}_{24})^{(4)}+G_{25}^0}{G_{25}^0}\right)}+(\hat{P}_{24})^{(4)}\right]$$

$(G_i^0)$ is as defined in the statement of theorem 4

The operator $\mathcal{A}^{(5)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that

$$G_{28}(t) \le G_{28}^0 + \int_0^t \left[(a_{28})^{(5)}\left(G_{29}^0+(\hat{P}_{28})^{(5)}e^{(\hat{M}_{28})^{(5)}s_{(28)}}\right)\right] ds_{(28)} =$$

$$\left(1+(a_{28})^{(5)}t\right)G_{29}^0 + \frac{(a_{28})^{(5)}(\hat{P}_{28})^{(5)}}{(\hat{M}_{28})^{(5)}}\left(e^{(\hat{M}_{28})^{(5)}t}-1\right)$$

From which it follows that 175

$$(G_{28}(t)-G_{28}^0)e^{-(\hat{M}_{28})^{(5)}t} \le \frac{(a_{28})^{(5)}}{(\hat{M}_{28})^{(5)}}\left[((\hat{P}_{28})^{(5)}+G_{29}^0)e^{\left(-\frac{(\hat{P}_{28})^{(5)}+G_{29}^0}{G_{29}^0}\right)}+(\hat{P}_{28})^{(5)}\right]$$

$(G_i^0)$ is as defined in the statement of theorem 5

The operator $\mathcal{A}^{(6)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that 176

$$G_{32}(t) \le G_{32}^0 + \int_0^t \left[(a_{32})^{(6)}\left(G_{33}^0+(\hat{P}_{32})^{(6)}e^{(\hat{M}_{32})^{(6)}s_{(32)}}\right)\right] ds_{(32)} =$$

$$\left(1+(a_{32})^{(6)}t\right)G_{33}^0 + \frac{(a_{32})^{(6)}(\hat{P}_{32})^{(6)}}{(\hat{M}_{32})^{(6)}}\left(e^{(\hat{M}_{32})^{(6)}t}-1\right)$$

From which it follows that 177

$$(G_{32}(t)-G_{32}^0)e^{-(\hat{M}_{32})^{(6)}t} \le \frac{(a_{32})^{(6)}}{(\hat{M}_{32})^{(6)}}\left[((\hat{P}_{32})^{(6)}+G_{33}^0)e^{\left(-\frac{(\hat{P}_{32})^{(6)}+G_{33}^0}{G_{33}^0}\right)}+(\hat{P}_{32})^{(6)}\right]$$

$(G_i^0)$ is as defined in the statement of theorem 6

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| Analogous inequalities hold also for $G_{25}, G_{26}, T_{24}, T_{25}, T_{26}$ | |
| (a) The operator $\mathcal{A}^{(7)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{36}(t) \le G_{36}^0 + \int_0^t \left[(a_{36})^{(7)} \left(G_{37}^0 + (\hat{P}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} s_{(36)}}\right)\right] ds_{(36)} =$<br>$\left(1 + (a_{36})^{(7)} t\right) G_{37}^0 + \frac{(a_{36})^{(7)} (\hat{P}_{36})^{(7)}}{(\hat{M}_{36})^{(7)}} \left(e^{(\hat{M}_{36})^{(7)} t} - 1\right)$ | 178 |
| From which it follows that<br>$(G_{36}(t) - G_{36}^0) e^{-(\hat{M}_{36})^{(7)} t} \le \frac{(a_{36})^{(7)}}{(\hat{M}_{36})^{(7)}} \left[\left((\hat{P}_{36})^{(7)} + G_{37}^0\right) e^{\left(-\frac{(\hat{P}_{36})^{(7)} + G_{37}^0}{G_{37}^0}\right)} + (\hat{P}_{36})^{(7)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 7 | |
| The operator $\mathcal{A}^{(8)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that | |
| $G_{40}(t) \le G_{40}^0 + \int_0^t \left[(a_{40})^{(8)} \left(G_{41}^0 + (\hat{P}_{40})^{(8)} e^{(\hat{M}_{40})^{(8)} s_{(40)}}\right)\right] ds_{(40)} =$<br>$\left(1 + (a_{40})^{(8)} t\right) G_{41}^0 + \frac{(a_{40})^{(8)} (\hat{P}_{40})^{(8)}}{(\hat{M}_{40})^{(8)}} \left(e^{(\hat{M}_{40})^{(8)} t} - 1\right)$ | 180 |
| From which it follows that<br>$(G_{40}(t) - G_{40}^0) e^{-(\hat{M}_{40})^{(8)} t} \le \frac{(a_{40})^{(8)}}{(\hat{M}_{40})^{(8)}} \left[\left((\hat{P}_{40})^{(8)} + G_{41}^0\right) e^{\left(-\frac{(\hat{P}_{40})^{(8)} + G_{41}^0}{G_{41}^0}\right)} + (\hat{P}_{40})^{(8)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 8<br>Analogous inequalities hold also for $G_{41}, G_{42}, T_{40}, T_{41}, T_{42}$ | 181 |
| The operator $\mathcal{A}^{(9)}$ maps the space of functions satisfying 34,35,36 into itself .Indeed it is obvious that<br>$G_{44}(t) \le G_{44}^0 + \int_0^t \left[(a_{44})^{(9)} \left(G_{45}^0 + (\hat{P}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} s_{(44)}}\right)\right] ds_{(44)} =$<br>$\left(1 + (a_{44})^{(9)} t\right) G_{45}^0 + \frac{(a_{44})^{(9)} (\hat{P}_{44})^{(9)}}{(\hat{M}_{44})^{(9)}} \left(e^{(\hat{M}_{44})^{(9)} t} - 1\right)$ | |
| From which it follows that<br>$(G_{44}(t) - G_{44}^0) e^{-(\hat{M}_{44})^{(9)} t} \le \frac{(a_{44})^{(9)}}{(\hat{M}_{44})^{(9)}} \left[\left((\hat{P}_{44})^{(9)} + G_{45}^0\right) e^{\left(-\frac{(\hat{P}_{44})^{(9)} + G_{45}^0}{G_{45}^0}\right)} + (\hat{P}_{44})^{(9)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 9<br>Analogous inequalities hold also for $G_{45}, G_{46}, T_{44}, T_{45}, T_{46}$ | |
| It is now sufficient to take $\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}} < 1$ and to choose | 182 |

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| $(\widehat{P}_{13})^{(1)}$ and $(\widehat{Q}_{13})^{(1)}$ large to have | |
| $\frac{(a_i)^{(1)}}{(\widehat{M}_{13})^{(1)}}\left[(\widehat{P}_{13})^{(1)}+\left((\hat{P}_{13})^{(1)}+G_j^0\right)e^{-\left(\frac{(\hat{P}_{13})^{(1)}+G_j^0}{G_j^0}\right)}\right]\leq(\hat{P}_{13})^{(1)}$ | 183 |
| $\frac{(b_i)^{(1)}}{(\widehat{M}_{13})^{(1)}}\left[\left((\hat{Q}_{13})^{(1)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{13})^{(1)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{13})^{(1)}\right]\leq(\hat{Q}_{13})^{(1)}$ | 184 |
| In order that the operator $\mathcal{A}^{(1)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(1)}$ is a contraction with respect to the metric<br>$d\left(\left(G^{(1)},T^{(1)}\right),\left(G^{(2)},T^{(2)}\right)\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{13})^{(1)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{13})^{(1)}t}\}$ | 185 |
| Indeed if we denote<br>**Definition of** $\tilde{G},\tilde{T}:\left(\tilde{G},\tilde{T}\right)=\mathcal{A}^{(1)}(G,T)$<br>It results<br>$\left\|\tilde{G}_{13}^{(1)}-\tilde{G}_i^{(2)}\right\|\leq\int_0^t(a_{13})^{(1)}\left\|G_{14}^{(1)}-G_{14}^{(2)}\right\|e^{-(\widehat{M}_{13})^{(1)}s_{(13)}}e^{(\widehat{M}_{13})^{(1)}s_{(13)}}ds_{(13)}+$<br>$\int_0^t\{(a_{13}')^{(1)}\left\|G_{13}^{(1)}-G_{13}^{(2)}\right\|e^{-(\widehat{M}_{13})^{(1)}s_{(13)}}e^{-(\widehat{M}_{13})^{(1)}s_{(13)}}+$<br>$(a_{13}'')^{(1)}\left(T_{14}^{(1)},s_{(13)}\right)\left\|G_{13}^{(1)}-G_{13}^{(2)}\right\|e^{-(\widehat{M}_{13})^{(1)}s_{(13)}}e^{(\widehat{M}_{13})^{(1)}s_{(13)}}+$<br>$G_{13}^{(2)}\left\\|(a_{13}'')^{(1)}\left(T_{14}^{(1)},s_{(13)}\right)-(a_{13}'')^{(1)}\left(T_{14}^{(2)},s_{(13)}\right)\right\\|\ e^{-(\widehat{M}_{13})^{(1)}s_{(13)}}e^{(\widehat{M}_{13})^{(1)}s_{(13)}}\}ds_{(13)}$<br>Where $s_{(13)}$ represents integrand that is integrated over the interval $[0,t]$<br>From the hypotheses it follows | |
| $\left\|G^{(1)}-G^{(2)}\right\|e^{-(\widehat{M}_{13})^{(1)}t}\leq$<br>$\frac{1}{(\widehat{M}_{13})^{(1)}}\left((a_{13})^{(1)}+(a_{13}')^{(1)}+(\widehat{A}_{13})^{(1)}+(\widehat{P}_{13})^{(1)}(\widehat{k}_{13})^{(1)}\right)d\left(\left(G^{(1)},T^{(1)};\ G^{(2)},T^{(2)}\right)\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 186 |
| **Remark 1:** The fact that we supposed $(a_{13}'')^{(1)}$ and $(b_{13}'')^{(1)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by | |

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| $(\widehat{P}_{13})^{(1)}e^{(\widehat{M}_{13})^{(1)}t}$ *and* $(\widehat{Q}_{13})^{(1)}e^{(\widehat{M}_{13})^{(1)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(1)}$ and $(b_i^{''})^{(1)}, i = 13,14,15$ depend only on $T_{14}$ and respectively on $G(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 2:** There does not exist any $t$ where $G_i(t) = 0\ and\ T_i(t) = 0$<br><br>From 19 to 24 it results<br><br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(1)} - (a_i^{''})^{(1)}(T_{14}(s_{(13)}), s_{(13)})\} ds_{(13)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(1)}t\right)} > 0$ for t > 0 | |
| **Definition of** $\left((\widehat{M}_{13})^{(1)}\right)_1, \left((\widehat{M}_{13})^{(1)}\right)_2\ and\ \left((\widehat{M}_{13})^{(1)}\right)_3$ :<br><br>**Remark 3**: if $G_{13}$ is bounded, the same property have also $G_{14}\ and\ G_{15}$ . indeed if<br><br>$G_{13} < (\widehat{M}_{13})^{(1)}$ it follows $\frac{dG_{14}}{dt} \leq \left((\widehat{M}_{13})^{(1)}\right)_1 - (a_{14}^{'})^{(1)}G_{14}$ and by integrating<br><br>$G_{14} \leq \left((\widehat{M}_{13})^{(1)}\right)_2 = G_{14}^0 + 2(a_{14})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_1/(a_{14}^{'})^{(1)}$<br><br>In the same way , one can obtain<br><br>$G_{15} \leq \left((\widehat{M}_{13})^{(1)}\right)_3 = G_{15}^0 + 2(a_{15})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_2/(a_{15}^{'})^{(1)}$<br><br>If $G_{14}\ or\ G_{15}$ is bounded, the same property follows for $G_{13}$ , $G_{15}$ and $G_{13}$ , $G_{14}$ respectively. | 187 |
| **Remark 4:** If $G_{13}$ *is* bounded, from below, the same property holds for $G_{14}\ and\ G_{15}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{14}$ is bounded from below. | 188 |
| **Remark 5:** If $T_{13}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(1)}(G(t),t)) = (b_{14}^{'})^{(1)}$ then $T_{14} \to \infty$.<br><br>**Definition of** $(m)^{(1)}$ and $\varepsilon_1$ :<br><br>Indeed let $t_1$ be so that for $t > t_1$<br><br>$(b_{14})^{(1)} - (b_i^{''})^{(1)}(G(t),t) < \varepsilon_1, T_{13}(t) > (m)^{(1)}$ | 189 |
| Then $\frac{dT_{14}}{dt} \geq (a_{14})^{(1)}(m)^{(1)} - \varepsilon_1 T_{14}$ which leads to<br><br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{\varepsilon_1}\right)(1 - e^{-\varepsilon_1 t}) + T_{14}^0 e^{-\varepsilon_1 t}$ If we take t such that $e^{-\varepsilon_1 t} = \frac{1}{2}$ it results<br><br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{2}\right),\quad t = log\frac{2}{\varepsilon_1}$ By taking now $\varepsilon_1$ sufficiently small one sees that $T_{14}$ is unbounded.<br>The same property holds for $T_{15}$ if $\lim_{t\to\infty}(b_{15}^{''})^{(1)}(G(t),t) = (b_{15}^{'})^{(1)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |

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| It is now sufficient to take $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}$ , $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ and to choose<br>$(\hat{P}_{16})^{(2)}$ *and* $(\hat{Q}_{16})^{(2)}$ large to have | 190 |
| $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[(\hat{P}_{16})^{(2)} + \left((\hat{P}_{16})^{(2)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{16})^{(2)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{16})^{(2)}$ | 191 |
| $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[\left((\hat{Q}_{16})^{(2)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{16})^{(2)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{16})^{(2)}\right] \leq (\hat{Q}_{16})^{(2)}$ | 192 |
| In order that the operator $\mathcal{A}^{(2)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 193 |
| The operator $\mathcal{A}^{(2)}$ is a contraction with respect to the metric<br>$d\left(((G_{19})^{(1)}, (T_{19})^{(1)}), ((G_{19})^{(2)}, (T_{19})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\hat{M}_{16})^{(2)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\hat{M}_{16})^{(2)}t}\}$ | 194 |
| Indeed if we denote<br>**Definition of** $\widetilde{G_{19}}, \widetilde{T_{19}}$ : $\left(\widetilde{G_{19}}, \widetilde{T_{19}}\right) = \mathcal{A}^{(2)}(G_{19}, T_{19})$ | 195 |
| It results<br>$\left\|\tilde{G}_{16}^{(1)} - \tilde{G}_i^{(2)}\right\| \leq \int_0^t (a_{16})^{(2)}\left\|G_{17}^{(1)} - G_{17}^{(2)}\right\|e^{-(\hat{M}_{16})^{(2)}s_{(16)}}e^{(\hat{M}_{16})^{(2)}s_{(16)}}\,ds_{(16)} +$<br>$\int_0^t\{(a_{16}')^{(2)}\left\|G_{16}^{(1)} - G_{16}^{(2)}\right\|e^{-(\hat{M}_{16})^{(2)}s_{(16)}}e^{-(\hat{M}_{16})^{(2)}s_{(16)}} +$<br>$(a_{16}'')^{(2)}\left(T_{17}^{(1)}, s_{(16)}\right)\left\|G_{16}^{(1)} - G_{16}^{(2)}\right\|e^{-(\hat{M}_{16})^{(2)}s_{(16)}}e^{(\hat{M}_{16})^{(2)}s_{(16)}} +$<br>$G_{16}^{(2)}\left\|(a_{16}'')^{(2)}\left(T_{17}^{(1)}, s_{(16)}\right) - (a_{16}'')^{(2)}\left(T_{17}^{(2)}, s_{(16)}\right)\right\| \; e^{-(\hat{M}_{16})^{(2)}s_{(16)}}e^{(\hat{M}_{16})^{(2)}s_{(16)}}\}ds_{(16)}$ | 196 |
| Where $s_{(16)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | 197 |
| $\left\|(G_{19})^{(1)} - (G_{19})^{(2)}\right\|e^{-(\hat{M}_{16})^{(2)}t} \leq$<br>$\frac{1}{(\hat{M}_{16})^{(2)}}\left((a_{16})^{(2)} + (a_{16}')^{(2)} + (\hat{A}_{16})^{(2)} + (\hat{P}_{16})^{(2)}(\hat{k}_{16})^{(2)}\right)d\left(((G_{19})^{(1)}, (T_{19})^{(1)}; (G_{19})^{(2)}, (T_{19})^{(2)})\right)$ | |
| And analogous inequalities for $G_i$ and $T_i$. Taking into account the hypothesis the result follows | 198 |
| **Remark 6:** The fact that we supposed $(a_{16}'')^{(2)}$ and $(b_{16}'')^{(2)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate | 199 |

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| condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{16})^{(2)}e^{(\widehat{M}_{16})^{(2)}t}$ and $(\widehat{Q}_{16})^{(2)}e^{(\widehat{M}_{16})^{(2)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}, i = 16,17,18$ depend only on $T_{17}$ and respectively on $(G_{19})$(and not on t) and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 7:** There does not exist any t where $G_i(t) = 0$ and $T_i(t) = 0$<br><br>it results<br><br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(2)} - (a_i^{''})^{(2)}(T_{17}(s_{(16)}), s_{(16)})\} ds_{(16)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{(-(b_i^{'})^{(2)}t)} > 0$ for $t > 0$ | 200 |
| **Definition of** $\left((\widehat{M}_{16})^{(2)}\right)_1, \left((\widehat{M}_{16})^{(2)}\right)_2$ and $\left((\widehat{M}_{16})^{(2)}\right)_3$ :<br><br>**Remark 8:** if $G_{16}$ is bounded, the same property have also $G_{17}$ and $G_{18}$ . indeed if<br><br>$G_{16} < (\widehat{M}_{16})^{(2)}$ it follows $\frac{dG_{17}}{dt} \leq \left((\widehat{M}_{16})^{(2)}\right)_1 - (a_{17}^{'})^{(2)}G_{17}$ and by integrating<br><br>$G_{17} \leq \left((\widehat{M}_{16})^{(2)}\right)_2 = G_{17}^0 + 2(a_{17})^{(2)}\left((\widehat{M}_{16})^{(2)}\right)_1/(a_{17}^{'})^{(2)}$<br><br>In the same way , one can obtain<br><br>$G_{18} \leq \left((\widehat{M}_{16})^{(2)}\right)_3 = G_{18}^0 + 2(a_{18})^{(2)}\left((\widehat{M}_{16})^{(2)}\right)_2/(a_{18}^{'})^{(2)}$<br><br>If $G_{17}$ or $G_{18}$ is bounded, the same property follows for $G_{16}$ , $G_{18}$ and $G_{16}$ , $G_{17}$ respectively. | 201 |
| **Remark 9**: If $G_{16}$ is bounded, from below, the same property holds for $G_{17}$ and $G_{18}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{17}$ is bounded from below. | 202 |
| **Remark 10:** If $T_{16}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(2)}((G_{19})(t), t)) = (b_{17}^{'})^{(2)}$ then $T_{17} \to \infty$.<br><br>**Definition of** $(m)^{(2)}$ and $\varepsilon_2$ :<br><br>Indeed let $t_2$ be so that for $t > t_2$<br><br>$(b_{17})^{(2)} - (b_i^{''})^{(2)}((G_{19})(t), t) < \varepsilon_2, T_{16}(t) > (m)^{(2)}$ | 203 |
| Then $\frac{dT_{17}}{dt} \geq (a_{17})^{(2)}(m)^{(2)} - \varepsilon_2 T_{17}$ which leads to<br><br>$T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{\varepsilon_2}\right)(1 - e^{-\varepsilon_2 t}) + T_{17}^0 e^{-\varepsilon_2 t}$ If we take t such that $e^{-\varepsilon_2 t} = \frac{1}{2}$ it results | 204 |
| $T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{2}\right)$, $t = \log\frac{2}{\varepsilon_2}$ By taking now $\varepsilon_2$ sufficiently small one sees that $T_{17}$ is unbounded.<br>The same property holds for $T_{18}$ if $\lim_{t\to\infty}(b_{18}^{''})^{(2)}\left((G_{19})(t), t\right) = (b_{18}^{'})^{(2)}$ | 205 |

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| We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |
| It is now sufficient to take $\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}$ , $\frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}} < 1$ and to choose<br>$(\hat{P}_{20})^{(3)}$ and $(\hat{Q}_{20})^{(3)}$ large to have | 207 |
| $\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}\left[(\hat{P}_{20})^{(3)} + \left((\hat{P}_{20})^{(3)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{20})^{(3)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{20})^{(3)}$ | 208 |
| $\frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}}\left[\left((\hat{Q}_{20})^{(3)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{20})^{(3)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{20})^{(3)}\right] \leq (\hat{Q}_{20})^{(3)}$ | 209 |
| In order that the operator $\mathcal{A}^{(3)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 210 |
| The operator $\mathcal{A}^{(3)}$ is a contraction with respect to the metric<br>$d\left(((G_{23})^{(1)}, (T_{23})^{(1)}), ((G_{23})^{(2)}, (T_{23})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\hat{M}_{20})^{(3)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\hat{M}_{20})^{(3)}t}\}$ | 211 |
| Indeed if we denote<br>**Definition of** $\widetilde{G_{23}}, \widetilde{T_{23}} :(\widetilde{(G_{23})}, \widetilde{(T_{23})}) = \mathcal{A}^{(3)}((G_{23}), (T_{23}))$ | 212 |
| It results<br>$\left\|\tilde{G}_{20}^{(1)} - \tilde{G}_i^{(2)}\right\| \leq \int_0^t (a_{20})^{(3)}\left\|G_{21}^{(1)} - G_{21}^{(2)}\right\|e^{-(\hat{M}_{20})^{(3)}s_{(20)}}e^{(\hat{M}_{20})^{(3)}s_{(20)}}\,ds_{(20)} +$<br>$\int_0^t\{(a'_{20})^{(3)}\left\|G_{20}^{(1)} - G_{20}^{(2)}\right\|e^{-(\hat{M}_{20})^{(3)}s_{(20)}}e^{-(\hat{M}_{20})^{(3)}s_{(20)}} +$<br>$(a''_{20})^{(3)}\left(T_{21}^{(1)}, s_{(20)}\right)\left\|G_{20}^{(1)} - G_{20}^{(2)}\right\|e^{-(\hat{M}_{20})^{(3)}s_{(20)}}e^{(\hat{M}_{20})^{(3)}s_{(20)}} +$<br>$G_{20}^{(2)}\left\|(a''_{20})^{(3)}\left(T_{21}^{(1)}, s_{(20)}\right) - (a''_{20})^{(3)}\left(T_{21}^{(2)}, s_{(20)}\right)\right\| \; e^{-(\hat{M}_{20})^{(3)}s_{(20)}}e^{(\hat{M}_{20})^{(3)}s_{(20)}}\}ds_{(20)}$<br>Where $s_{(20)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | 213 |
| $\left\|G_{23}{}^{(1)} - G_{23}{}^{(2)}\right\|e^{-(\hat{M}_{20})^{(3)}t} \leq$<br>$\frac{1}{(\hat{M}_{20})^{(3)}}\left((a_{20})^{(3)} + (a'_{20})^{(3)} + (\hat{A}_{20})^{(3)} + (\hat{P}_{20})^{(3)}(\hat{k}_{20})^{(3)}\right)d\left(((G_{23})^{(1)}, (T_{23})^{(1)};\ (G_{23})^{(2)}, (T_{23})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 214 |

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| **Remark 11:** The fact that we supposed $(a_{20}^{''})^{(3)}$ and $(b_{20}^{''})^{(3)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}$ $and$ $(\widehat{Q}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(3)}$ and $(b_i^{''})^{(3)}, i = 20,21,22$ depend only on $T_{21}$ and respectively on $(G_{23})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 215 |
| **Remark 12:** There does not exist any $t$ where $G_i(t) = 0\ and\ T_i(t) = 0$<br><br>it results<br><br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(3)} - (a_i^{''})^{(3)}(T_{21}(s_{(20)}), s_{(20)})\} ds_{(20)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{(-(b_i^{'})^{(3)}t)} > 0$ for t > 0 | 216 |
| **Definition of** $\left((\widehat{M}_{20})^{(3)}\right)_1, \left((\widehat{M}_{20})^{(3)}\right)_2\ and\ \left((\widehat{M}_{20})^{(3)}\right)_3$ :<br><br>**Remark 13:** if $G_{20}$ is bounded, the same property have also $G_{21}\ and\ G_{22}$ . indeed if<br><br>$G_{20} < (\widehat{M}_{20})^{(3)}$ it follows $\frac{dG_{21}}{dt} \leq \left((\widehat{M}_{20})^{(3)}\right)_1 - (a_{21}^{'})^{(3)}G_{21}$ and by integrating<br><br>$G_{21} \leq \left((\widehat{M}_{20})^{(3)}\right)_2 = G_{21}^0 + 2(a_{21})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_1/(a_{21}^{'})^{(3)}$<br><br>In the same way , one can obtain<br><br>$G_{22} \leq \left((\widehat{M}_{20})^{(3)}\right)_3 = G_{22}^0 + 2(a_{22})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_2/(a_{22}^{'})^{(3)}$<br><br>If $G_{21}\ or\ G_{22}$ is bounded, the same property follows for $G_{20}$ , $G_{22}$ and $G_{20}$ , $G_{21}$ respectively. | 217 |
| **Remark 14:** If $G_{20}$ $is$ bounded, from below, the same property holds for $G_{21}\ and\ G_{22}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{21}$ is bounded from below. | 218 |
| **Remark 15:** If $T_{20}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(3)}\left((G_{23})(t), t\right)) = (b_{21}^{'})^{(3)}$ then $T_{21} \to \infty$.<br><br>**Definition of** $(m)^{(3)}$ and $\varepsilon_3$ :<br><br>Indeed let $t_3$ be so that for $t > t_3$<br><br>$(b_{21})^{(3)} - (b_i^{''})^{(3)}\left((G_{23})(t), t\right) < \varepsilon_3, T_{20}(t) > (m)^{(3)}$ | 219 |
| Then $\frac{dT_{21}}{dt} \geq (a_{21})^{(3)}(m)^{(3)} - \varepsilon_3 T_{21}$ which leads to<br><br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{\varepsilon_3}\right)(1 - e^{-\varepsilon_3 t}) + T_{21}^0 e^{-\varepsilon_3 t}$ If we take t such that $e^{-\varepsilon_3 t} = \frac{1}{2}$ it results<br><br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{2}\right)$, $t = log\frac{2}{\varepsilon_3}$ By taking now $\varepsilon_3$ sufficiently small one sees that $T_{21}$ is unbounded. | 220 |

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| The same property holds for $T_{22}$ if $\lim_{t\to\infty}(b_{22}'')^{(3)}\left((G_{23})(t),t\right)=(b_{22}')^{(3)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |
| It is now sufficient to take $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}$ , $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}<1$ and to choose<br>$(\widehat{P}_{24})^{(4)}$ and $(\widehat{Q}_{24})^{(4)}$ large to have | 221 |
| $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[(\widehat{P}_{24})^{(4)}+\left((\hat{P}_{24})^{(4)}+G_j^0\right)e^{-\left(\frac{(\hat{P}_{24})^{(4)}+G_j^0}{G_j^0}\right)}\right]\le(\hat{P}_{24})^{(4)}$ | 222 |
| $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[\left((\hat{Q}_{24})^{(4)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{24})^{(4)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{24})^{(4)}\right]\le(\hat{Q}_{24})^{(4)}$ | 223 |
| In order that the operator $\mathcal{A}^{(4)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 224 |
| The operator $\mathcal{A}^{(4)}$ is a contraction with respect to the metric<br>$d\left(((G_{27})^{(1)},(T_{27})^{(1)}),((G_{27})^{(2)},(T_{27})^{(2)})\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{24})^{(4)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{24})^{(4)}t}\}$<br>Indeed if we denote<br>**Definition of** $(\widetilde{G_{27}}),(\widetilde{T_{27}})$ : $\left((\widetilde{G_{27}}),(\widetilde{T_{27}})\right)=\mathcal{A}^{(4)}((G_{27}),(T_{27}))$<br>It results<br>$\left\|\tilde{G}_{24}^{(1)}-\tilde{G}_i^{(2)}\right\|\le\int_0^t(a_{24})^{(4)}\left\|G_{25}^{(1)}-G_{25}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}ds_{(24)}+$<br>$\int_0^t\{(a_{24}')^{(4)}\left\|G_{24}^{(1)}-G_{24}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}+$<br>$(a_{24}'')^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)\left\|G_{24}^{(1)}-G_{24}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}+$<br>$G_{24}^{(2)}\left\\|(a_{24}'')^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)-(a_{24}'')^{(4)}\left(T_{25}^{(2)},s_{(24)}\right)\right\\|\ e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}\}ds_{(24)}$<br>Where $s_{(24)}$ represents integrand that is integrated over the interval $[0,t]$<br>From the hypotheses on Equations it follows | 225 |
| $\left\|(G_{27})^{(1)}-(G_{27})^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}t}\le$<br>$\frac{1}{(\widehat{M}_{24})^{(4)}}\left((a_{24})^{(4)}+(a_{24}')^{(4)}+(\widehat{A}_{24})^{(4)}+(\widehat{P}_{24})^{(4)}(\widehat{k}_{24})^{(4)}\right)d\left(((G_{27})^{(1)},(T_{27})^{(1)};\ (G_{27})^{(2)},(T_{27})^{(2)})\right)$ | 226 |

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| And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | |
| **Remark 16:** The fact that we supposed $(a_{24}^{''})^{(4)}$ and $(b_{24}^{''})^{(4)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}$ *and* $(\widehat{Q}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(4)}$ and $(b_i^{''})^{(4)}, i = 24,25,26$ depend only on $T_{25}$ and respectively on $(G_{27})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 227 |
| **Remark 17:** There does not exist any $t$ where $G_i\,(t) = 0\ and\ T_i\,(t) = 0$<br>it results<br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t\{(a_i^{'})^{(4)}-(a_i^{''})^{(4)}(T_{25}(s_{(24)}),s_{(24)})\}ds_{(24)}\right]} \geq 0$<br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(4)}t\right)} > 0$ for t > 0 | 228 |
| **Definition of** $\left((\widehat{M}_{24})^{(4)}\right)_1, \left((\widehat{M}_{24})^{(4)}\right)_2\ and\ \left((\widehat{M}_{24})^{(4)}\right)_3$ :<br>**Remark 18:** if $G_{24}$ is bounded, the same property have also $G_{25}\ and\ G_{26}$ . indeed if<br>$G_{24} < (\widehat{M}_{24})^{(4)}$ it follows $\frac{dG_{25}}{dt} \leq \left((\widehat{M}_{24})^{(4)}\right)_1 - (a_{25}^{'})^{(4)}G_{25}$ and by integrating<br>$G_{25} \leq \left((\widehat{M}_{24})^{(4)}\right)_2 = G_{25}^0 + 2(a_{25})^{(4)}\left((\widehat{M}_{24})^{(4)}\right)_1/(a_{25}^{'})^{(4)}$<br>In the same way , one can obtain<br>$G_{26} \leq \left((\widehat{M}_{24})^{(4)}\right)_3 = G_{26}^0 + 2(a_{26})^{(4)}\left((\widehat{M}_{24})^{(4)}\right)_2/(a_{26}^{'})^{(4)}$<br>If $G_{25}\ or\ G_{26}$ is bounded, the same property follows for $G_{24}$ , $G_{26}$ and $G_{24}$ , $G_{25}$ respectively. | 229 |
| **Remark 19:** If $G_{24}$ *is* bounded, from below, the same property holds for $G_{25}\ and\ G_{26}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{25}$ is bounded from below. | 230 |
| **Remark 20:** If $T_{24}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(4)}((G_{27})(t),t)) = (b_{25}^{'})^{(4)}$ then $T_{25} \to \infty$.<br>**Definition of** $(m)^{(4)}$ and $\varepsilon_4$ :<br>Indeed let $t_4$ be so that for $t > t_4$<br>$(b_{25})^{(4)} - (b_i^{''})^{(4)}((G_{27})(t),t) < \varepsilon_4, T_{24}\,(t) > (m)^{(4)}$ | 231 |
| Then $\frac{dT_{25}}{dt} \geq (a_{25})^{(4)}(m)^{(4)} - \varepsilon_4 T_{25}$ which leads to<br>$T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{\varepsilon_4}\right)(1 - e^{-\varepsilon_4 t}) + T_{25}^0 e^{-\varepsilon_4 t}$ If we take t such that $e^{-\varepsilon_4 t} = \frac{1}{2}$ it results | 232 |

$T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{2}\right), \quad t = log\frac{2}{\varepsilon_4}$ By taking now $\varepsilon_4$ sufficiently small one sees that $T_{25}$ is unbounded. The same property holds for $T_{26}$ if $\lim_{t\to\infty}(b_{26}^{''})^{(4)}\left((G_{27})(t),t\right) = (b_{26}^{'})^{(4)}$

We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 42

Analogous inequalities hold also for $G_{29}, G_{30}, T_{28}, T_{29}, T_{30}$

It is now sufficient to take $\frac{(a_i)^{(5)}}{(\hat{M}_{28})^{(5)}}, \frac{(b_i)^{(5)}}{(\hat{M}_{28})^{(5)}} < 1$ and to choose

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$(\hat{P}_{28})^{(5)}$ and $(\hat{Q}_{28})^{(5)}$ large to have

$$\frac{(a_i)^{(5)}}{(\hat{M}_{28})^{(5)}}\left[(\hat{P}_{28})^{(5)} + \left((\hat{P}_{28})^{(5)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{28})^{(5)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{28})^{(5)} \qquad 234$$

$$\frac{(b_i)^{(5)}}{(\hat{M}_{28})^{(5)}}\left[\left((\hat{Q}_{28})^{(5)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{28})^{(5)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{28})^{(5)}\right] \leq (\hat{Q}_{28})^{(5)} \qquad 235$$

In order that the operator $\mathcal{A}^{(5)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself

The operator $\mathcal{A}^{(5)}$ is a contraction with respect to the metric 236

$$d\left(((G_{31})^{(1)},(T_{31})^{(1)}),((G_{31})^{(2)},(T_{31})^{(2)})\right) =$$

$$\sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t) - G_i^{(2)}(t)\right|e^{-(\hat{M}_{28})^{(5)}t}, \max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t) - T_i^{(2)}(t)\right|e^{-(\hat{M}_{28})^{(5)}t}\}$$

Indeed if we denote

**Definition of** $(\widetilde{G_{31}}),(\widetilde{T_{31}})$ : $\left((\widetilde{G_{31}}),(\widetilde{T_{31}})\right) = \mathcal{A}^{(5)}\left((G_{31}),(T_{31})\right)$

It results

$$\left|\tilde{G}_{28}^{(1)} - \tilde{G}_i^{(2)}\right| \leq \int_0^t (a_{28})^{(5)}\left|G_{29}^{(1)} - G_{29}^{(2)}\right|e^{-(\hat{M}_{28})^{(5)}s_{(28)}}e^{(\hat{M}_{28})^{(5)}s_{(28)}}\,ds_{(28)} +$$

$$\int_0^t\{(a_{28}^{'})^{(5)}\left|G_{28}^{(1)} - G_{28}^{(2)}\right|e^{-(\hat{M}_{28})^{(5)}s_{(28)}}e^{-(\hat{M}_{28})^{(5)}s_{(28)}} +$$

$$(a_{28}^{''})^{(5)}\left(T_{29}^{(1)},s_{(28)}\right)\left|G_{28}^{(1)} - G_{28}^{(2)}\right|e^{-(\hat{M}_{28})^{(5)}s_{(28)}}e^{(\hat{M}_{28})^{(5)}s_{(28)}} +$$

$$G_{28}^{(2)}\left|(a_{28}^{''})^{(5)}\left(T_{29}^{(1)},s_{(28)}\right) - (a_{28}^{''})^{(5)}\left(T_{29}^{(2)},s_{(28)}\right)\right|\ e^{-(\hat{M}_{28})^{(5)}s_{(28)}}e^{(\hat{M}_{28})^{(5)}s_{(28)}}\}ds_{(28)}$$

Where $s_{(28)}$ represents integrand that is integrated over the interval [0, t]

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| From the hypotheses on it follows | |
| $\left\|(G_{31})^{(1)} - (G_{31})^{(2)}\right\| e^{-(\widehat{M}_{28})^{(5)}t} \leq$ $\frac{1}{(\widehat{M}_{28})^{(5)}}\left((a_{28})^{(5)} + (a'_{28})^{(5)} + (\widehat{A}_{28})^{(5)} + (\widehat{P}_{28})^{(5)}(\widehat{k}_{28})^{(5)}\right) d\left(((G_{31})^{(1)}, (T_{31})^{(1)};\ (G_{31})^{(2)}, (T_{31})^{(2)})\right)$ And analogous inequalities for $G_i$ and $T_i$. Taking into account the hypothesis the result follows | 237 |
| **Remark 21:** The fact that we supposed $(a''_{28})^{(5)}$ and $(b''_{28})^{(5)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{28})^{(5)} e^{(\widehat{M}_{28})^{(5)}t}$ and $(\widehat{Q}_{28})^{(5)} e^{(\widehat{M}_{28})^{(5)}t}$ respectively of $\mathbb{R}_+$. If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a''_i)^{(5)}$ and $(b''_i)^{(5)}, i = 28,29,30$ depend only on $T_{29}$ and respectively on $(G_{31})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 238 |
| **Remark 22:** There does not exist any $t$ where $G_i(t) = 0$ and $T_i(t) = 0$ it results $G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a'_i)^{(5)} - (a''_i)^{(5)}(T_{29}(s_{(28)}), s_{(28)})\} ds_{(28)}\right]} \geq 0$ $T_i(t) \geq T_i^0 e^{\left(-(b'_i)^{(5)}t\right)} > 0$ for t > 0 | 239 |
| **Definition of** $\left((\widehat{M}_{28})^{(5)}\right)_1, \left((\widehat{M}_{28})^{(5)}\right)_2$ and $\left((\widehat{M}_{28})^{(5)}\right)_3$ : **Remark 23:** if $G_{28}$ is bounded, the same property have also $G_{29}$ and $G_{30}$ . indeed if $G_{28} < (\widehat{M}_{28})^{(5)}$ it follows $\frac{dG_{29}}{dt} \leq \left((\widehat{M}_{28})^{(5)}\right)_1 - (a'_{29})^{(5)} G_{29}$ and by integrating $G_{29} \leq \left((\widehat{M}_{28})^{(5)}\right)_2 = G_{29}^0 + 2(a_{29})^{(5)}\left((\widehat{M}_{28})^{(5)}\right)_1 / (a'_{29})^{(5)}$ In the same way , one can obtain $G_{30} \leq \left((\widehat{M}_{28})^{(5)}\right)_3 = G_{30}^0 + 2(a_{30})^{(5)}\left((\widehat{M}_{28})^{(5)}\right)_2 / (a'_{30})^{(5)}$ If $G_{29}$ or $G_{30}$ is bounded, the same property follows for $G_{28}$ , $G_{30}$ and $G_{28}$ , $G_{29}$ respectively. | 240 |
| **Remark 24:** If $G_{28}$ is bounded, from below, the same property holds for $G_{29}$ and $G_{30}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{29}$ is bounded from below. | 241 |

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| **Remark 25:** If $T_{28}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(5)}((G_{31})(t),t)) = (b_{29}^{'})^{(5)}$ then $T_{29} \to \infty$.<br><br>**Definition of** $(m)^{(5)}$ and $\varepsilon_5$ :<br><br>Indeed let $t_5$ be so that for $t > t_5$<br><br>$(b_{29})^{(5)} - (b_i^{''})^{(5)}((G_{31})(t),t) < \varepsilon_5, T_{28}(t) > (m)^{(5)}$ | 242 |
| Then $\frac{dT_{29}}{dt} \geq (a_{29})^{(5)}(m)^{(5)} - \varepsilon_5 T_{29}$ which leads to<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{\varepsilon_5}\right)(1-e^{-\varepsilon_5 t}) + T_{29}^0 e^{-\varepsilon_5 t}$ If we take t such that $e^{-\varepsilon_5 t} = \frac{1}{2}$ it results<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{2}\right), \quad t = log\frac{2}{\varepsilon_5}$ By taking now $\varepsilon_5$ sufficiently small one sees that $T_{29}$ is unbounded.<br>The same property holds for $T_{30}$ if $\lim_{t\to\infty}(b_{30}^{''})^{(5)}\left((G_{31})(t),t\right) = (b_{30}^{'})^{(5)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations<br><br>Analogous inequalities hold also for $G_{33}, G_{34}, T_{32}, T_{33}, T_{34}$ | 243 |
| It is now sufficient to take $\frac{(a_i)^{(6)}}{(\hat{M}_{32})^{(6)}}, \frac{(b_i)^{(6)}}{(\hat{M}_{32})^{(6)}} < 1$ and to choose<br><br>$(\hat{P}_{32})^{(6)}$ and $(\hat{Q}_{32})^{(6)}$ large to have | 244 |
| $\frac{(a_i)^{(6)}}{(\hat{M}_{32})^{(6)}}\left[(\hat{P}_{32})^{(6)} + \left((\hat{P}_{32})^{(6)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{32})^{(6)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{32})^{(6)}$ | 245 |
| $\frac{(b_i)^{(6)}}{(\hat{M}_{32})^{(6)}}\left[\left((\hat{Q}_{32})^{(6)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{32})^{(6)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{32})^{(6)}\right] \leq (\hat{Q}_{32})^{(6)}$ | 246 |
| In order that the operator $\mathcal{A}^{(6)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(6)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{35})^{(1)},(T_{35})^{(1)}),((G_{35})^{(2)},(T_{35})^{(2)})\right) =$<br><br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t) - G_i^{(2)}(t)\right|e^{-(\hat{M}_{32})^{(6)}t}, \max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t) - T_i^{(2)}(t)\right|e^{-(\hat{M}_{32})^{(6)}t}\}$<br><br>Indeed if we denote<br><br>**Definition of** $\widetilde{(G_{35})}, \widetilde{(T_{35})}$ : $\left(\widetilde{(G_{35})}, \widetilde{(T_{35})}\right) = \mathcal{A}^{(6)}\left((G_{35}),(T_{35})\right)$<br><br>It results | 247 |

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| $\left\|\tilde{G}_{32}^{(1)}-\tilde{G}_{i}^{(2)}\right\| \leq \int_0^t (a_{32})^{(6)} \left\|G_{33}^{(1)}-G_{33}^{(2)}\right\| e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} e^{(\widehat{M}_{32})^{(6)}s_{(32)}}\, ds_{(32)} +$<br>$\int_0^t \{(a_{32}^{'})^{(6)} \left\|G_{32}^{(1)}-G_{32}^{(2)}\right\| e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} +$<br>$(a_{32}^{''})^{(6)}\left(T_{33}^{(1)}, s_{(32)}\right)\left\|G_{32}^{(1)}-G_{32}^{(2)}\right\| e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} e^{(\widehat{M}_{32})^{(6)}s_{(32)}} +$<br>$G_{32}^{(2)} \left\|(a_{32}^{''})^{(6)}\left(T_{33}^{(1)}, s_{(32)}\right) - (a_{32}^{''})^{(6)}\left(T_{33}^{(2)}, s_{(32)}\right)\right\| \; e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} e^{(\widehat{M}_{32})^{(6)}s_{(32)}}\} ds_{(32)}$<br>Where $s_{(32)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | |
| $\left\|(G_{35})^{(1)}-(G_{35})^{(2)}\right\| e^{-(\widehat{M}_{32})^{(6)}t} \leq$<br>$\frac{1}{(\widehat{M}_{32})^{(6)}}\left((a_{32})^{(6)} + (a_{32}^{'})^{(6)} + (\widehat{A}_{32})^{(6)} + (\widehat{P}_{32})^{(6)}(\widehat{k}_{32})^{(6)}\right) d\left(((G_{35})^{(1)}, (T_{35})^{(1)};\ (G_{35})^{(2)}, (T_{35})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 248 |
| **Remark 26:** The fact that we supposed $(a_{32}^{''})^{(6)}$ and $(b_{32}^{''})^{(6)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{32})^{(6)} e^{(\widehat{M}_{32})^{(6)}t}$ *and* $(\widehat{Q}_{32})^{(6)} e^{(\widehat{M}_{32})^{(6)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(6)}$ and $(b_i^{''})^{(6)}, i = 32,33,34$ depend only on $T_{33}$ and respectively on $(G_{35})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 249 |
| **Remark 27:** There does not exist any $t$ where $G_i\,(t) = 0$ *and* $T_i\,(t) = 0$<br>it results<br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(6)} - (a_i^{''})^{(6)}(T_{33}(s_{(32)}), s_{(32)})\} ds_{(32)}\right]} \geq 0$<br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(6)}t\right)} > 0$ for t > 0 | 250 |
| **Definition of** $\left((\widehat{M}_{32})^{(6)}\right)_1, \left((\widehat{M}_{32})^{(6)}\right)_2$ *and* $\left((\widehat{M}_{32})^{(6)}\right)_3$ :<br>**Remark 28:** if $G_{32}$ is bounded, the same property have also $G_{33}$ *and* $G_{34}$ . indeed if<br>$G_{32} < (\widehat{M}_{32})^{(6)}$ it follows $\frac{dG_{33}}{dt} \leq \left((\widehat{M}_{32})^{(6)}\right)_1 - (a_{33}^{'})^{(6)} G_{33}$ and by integrating<br>$G_{33} \leq \left((\widehat{M}_{32})^{(6)}\right)_2 = G_{33}^0 + 2(a_{33})^{(6)}\left((\widehat{M}_{32})^{(6)}\right)_1 / (a_{33}^{'})^{(6)}$<br>In the same way , one can obtain<br>$G_{34} \leq \left((\widehat{M}_{32})^{(6)}\right)_3 = G_{34}^0 + 2(a_{34})^{(6)}\left((\widehat{M}_{32})^{(6)}\right)_2 / (a_{34}^{'})^{(6)}$ | 251 |

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| If $G_{33}$ *or* $G_{34}$ is bounded, the same property follows for $G_{32}$ , $G_{34}$ and $G_{32}$ , $G_{33}$ respectively. | |
| **Remark 29:** If $G_{32}$ *is* bounded, from below, the same property holds for $G_{33}$ *and* $G_{34}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{33}$ is bounded from below. | 252 |
| **Remark 30:** If $T_{32}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(6)}((G_{35})(t),t)) = (b_{33}^{'})^{(6)}$ then $T_{33} \to \infty$.<br><br>**Definition of** $(m)^{(6)}$ and $\varepsilon_6$ :<br><br>Indeed let $t_6$ be so that for $t > t_6$<br><br>$(b_{33})^{(6)} - (b_i^{''})^{(6)}\big((G_{35})(t),t\big) < \varepsilon_6, T_{32}(t) > (m)^{(6)}$ | 253 |
| Then $\frac{dT_{33}}{dt} \geq (a_{33})^{(6)}(m)^{(6)} - \varepsilon_6 T_{33}$ which leads to<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{\varepsilon_6}\right)(1 - e^{-\varepsilon_6 t}) + T_{33}^0 e^{-\varepsilon_6 t}$ If we take t such that $e^{-\varepsilon_6 t} = \frac{1}{2}$ it results<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{2}\right)$, $t = log\frac{2}{\varepsilon_6}$ By taking now $\varepsilon_6$ sufficiently small one sees that $T_{33}$ is unbounded.<br>The same property holds for $T_{34}$ if $\lim_{t\to\infty}(b_{34}^{''})^{(6)}\left((G_{35})(t),t(t),t\right) = (b_{34}^{'})^{(6)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 254 |
| Analogous inequalities hold also for $G_{37}$ , $G_{38}, T_{36}, T_{37}, T_{38}$<br><br>It is now sufficient to take $\frac{(a_i)^{(7)}}{(\hat{M}_{36})^{(7)}}$ , $\frac{(b_i)^{(7)}}{(\hat{M}_{36})^{(7)}} < 1$ and to choose<br>$(\hat{P}_{36})^{(7)}$ and $(\hat{Q}_{36})^{(7)}$ large to have | 255 |
| $\frac{(a_i)^{(7)}}{(\hat{M}_{36})^{(7)}}\left[(\hat{P}_{36})^{(7)} + \left((\hat{P}_{36})^{(7)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{36})^{(7)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{36})^{(7)}$ | 256 |
| $\frac{(b_i)^{(7)}}{(\hat{M}_{36})^{(7)}}\left[\left((\hat{Q}_{36})^{(7)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{36})^{(7)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{36})^{(7)}\right] \leq (\hat{Q}_{36})^{(7)}$ | 257 |
| In order that the operator $\mathcal{A}^{(7)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(7)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{39})^{(1)},(T_{39})^{(1)}),((G_{39})^{(2)},(T_{39})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t) - G_i^{(2)}(t)\right|e^{-(\hat{M}_{36})^{(7)}t}, \max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t) - T_i^{(2)}(t)\right|e^{-(\hat{M}_{36})^{(7)}t}\}$<br><br>Indeed if we denote | 258 |

**Definition of** $\widetilde{(G_{39})}, \widetilde{(T_{39})} : \left( \widetilde{(G_{39})}, \widetilde{(T_{39})} \right) = \mathcal{A}^{(7)}((G_{39}), (T_{39}))$
It results

$\left|\tilde{G}_{36}^{(1)} - \tilde{G}_{i}^{(2)}\right| \leq \int_0^t (a_{36})^{(7)} \left|G_{37}^{(1)} - G_{37}^{(2)}\right| e^{-(\widehat{M}_{36})^{(7)} s_{(36)}} e^{(\widehat{M}_{36})^{(7)} s_{(36)}}\, ds_{(36)} +$
$\int_0^t \{(a_{36}^{'})^{(7)} \left|G_{36}^{(1)} - G_{36}^{(2)}\right| e^{-(\widehat{M}_{36})^{(7)} s_{(36)}} e^{-(\widehat{M}_{36})^{(7)} s_{(36)}} +$
$(a_{36}^{''})^{(7)}\left(T_{37}^{(1)}, s_{(36)}\right) \left|G_{36}^{(1)} - G_{36}^{(2)}\right| e^{-(\widehat{M}_{36})^{(7)} s_{(36)}} e^{(\widehat{M}_{36})^{(7)} s_{(36)}} +$
$G_{36}^{(2)} \left|(a_{36}^{''})^{(7)}\left(T_{37}^{(1)}, s_{(36)}\right) - (a_{36}^{''})^{(7)}\left(T_{37}^{(2)}, s_{(36)}\right)\right| \; e^{-(\widehat{M}_{36})^{(7)} s_{(36)}} e^{(\widehat{M}_{36})^{(7)} s_{(36)}} \} ds_{(36)}$

Where $s_{(36)}$ represents integrand that is integrated over the interval $[0, t]$

From the hypotheses on it follows

$\left|(G_{39})^{(1)} - (G_{39})^{(2)}\right| e^{-(\widehat{M}_{36})^{(7)} t} \leq$
$\frac{1}{(\widehat{M}_{36})^{(7)}} \left((a_{36})^{(7)} + (a_{36}^{'})^{(7)} + (\widehat{A}_{36})^{(7)} + (\widehat{P}_{36})^{(7)} (\widehat{k}_{36})^{(7)}\right) d\left(((G_{39})^{(1)}, (T_{39})^{(1)};\ (G_{39})^{(2)}, (T_{39})^{(2)})\right)$

And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows

**Remark 31:** The fact that we supposed $(a_{36}^{''})^{(7)}$ and $(b_{36}^{''})^{(7)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{36})^{(7)} e^{(\widehat{M}_{36})^{(7)} t}$ *and* $(\widehat{Q}_{36})^{(7)} e^{(\widehat{M}_{36})^{(7)} t}$ respectively of $\mathbb{R}_+$.

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(7)}$ and $(b_i^{''})^{(7)}, i = 36{,}37{,}38$ depend only on $T_{37}$ and respectively on $(G_{39})$(*and not on* $t$) and hypothesis can replaced by a usual Lipschitz condition.

**Remark 32:** There does not exist any $t$ where $G_i\,(t) = 0$ *and* $T_i\,(t) = 0$

it results

$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(7)} - (a_i^{''})^{(7)} (T_{37}(s_{(36)}), s_{(36)})\} ds_{(36)}\right]} \geq 0$

$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(7)} t\right)} > 0$ for t > 0

**Definition of** $\left((\widehat{M}_{36})^{(7)}\right)_1, \left((\widehat{M}_{36})^{(7)}\right)_2$ *and* $\left((\widehat{M}_{36})^{(7)}\right)_3$ :

**Remark 33:** if $G_{36}$ is bounded, the same property have also $G_{37}$ *and* $G_{38}$ . indeed if

$G_{36} < (\widehat{M}_{36})^{(7)}$ it follows $\frac{dG_{37}}{dt} \leq \left((\widehat{M}_{36})^{(7)}\right)_1 - (a_{37}^{'})^{(7)} G_{37}$ and by integrating

$G_{37} \leq \left((\widehat{M}_{36})^{(7)}\right)_2 = G_{37}^0 + 2(a_{37})^{(7)} \left((\widehat{M}_{36})^{(7)}\right)_1 / (a_{37}^{'})^{(7)}$

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| In the same way , one can obtain<br>$G_{38} \leq \left((\widehat{M}_{36})^{(7)}\right)_3 = G_{38}^0 + 2(a_{38})^{(7)}\left((\widehat{M}_{36})^{(7)}\right)_2/(a_{38}')^{(7)}$<br>If $G_{37}$ *or* $G_{38}$ is bounded, the same property follows for $G_{36}$ , $G_{38}$ and $G_{36}$ , $G_{37}$ respectively. | |
| **Remark 34:** If $G_{36}$ *is* bounded, from below, the same property holds for $G_{37}$ *and* $G_{38}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{37}$ is bounded from below. | 263 |
| **Remark 35:** If $T_{36}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(7)}((G_{39})(t),t)) = (b_{37}')^{(7)}$ then $T_{37} \to \infty$.<br>**Definition of** $(m)^{(7)}$ and $\varepsilon_7$ :<br>Indeed let $t_7$ be so that for $t > t_7$<br>$(b_{37})^{(7)} - (b_i'')^{(7)}((G_{39})(t),t) < \varepsilon_7, T_{36}(t) > (m)^{(7)}$ | 264 |
| Then $\frac{dT_{37}}{dt} \geq (a_{37})^{(7)}(m)^{(7)} - \varepsilon_7 T_{37}$ which leads to<br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{\varepsilon_7}\right)(1 - e^{-\varepsilon_7 t}) + T_{37}^0 e^{-\varepsilon_7 t}$ If we take t such that $e^{-\varepsilon_7 t} = \frac{1}{2}$ it results<br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{2}\right)$, $t = log\frac{2}{\varepsilon_7}$ By taking now $\varepsilon_7$ sufficiently small one sees that $T_{37}$ is unbounded.<br>The same property holds for $T_{38}$ if $\lim_{t\to\infty}(b_{38}'')^{(7)}\left((G_{39})(t),t\right) = (b_{38}')^{(7)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 265 |
| It is now sufficient to take $\frac{(a_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}$ , $\frac{(b_i)^{(8)}}{(\widehat{M}_{40})^{(8)}} < 1$ and to choose<br>$(\widehat{P}_{40})^{(8)}$ and $(\widehat{Q}_{40})^{(8)}$ large to have | 266 |
| $\frac{(a_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}\left[(\widehat{P}_{40})^{(8)} + \left((\widehat{P}_{40})^{(8)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{40})^{(8)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{40})^{(8)}$ | 267 |
| $\frac{(b_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}\left[\left((\widehat{Q}_{40})^{(8)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{40})^{(8)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{40})^{(8)}\right] \leq (\widehat{Q}_{40})^{(8)}$ | 268 |
| In order that the operator $\mathcal{A}^{(8)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(8)}$ is a contraction with respect to the metric | |

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| $d\left(((G_{43})^{(1)},(T_{43})^{(1)}),((G_{43})^{(2)},(T_{43})^{(2)})\right) =$ $\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{40})^{(8)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{40})^{(8)}t}\}$ | 269 |
| Indeed if we denote<br><br>**Definition of** $\widetilde{(G_{43})},\widetilde{(T_{43})}$ : $\left(\widetilde{(G_{43})},\widetilde{(T_{43})}\right)=\mathcal{A}^{(8)}((G_{43}),(T_{43}))$ | 270 |
| It results<br><br>$\left\|\tilde{G}_{40}^{(1)}-\tilde{G}_i^{(2)}\right\|\le\int_0^t(a_{40})^{(8)}\left\|G_{41}^{(1)}-G_{41}^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{(\widehat{M}_{40})^{(8)}s_{(40)}}\,ds_{(40)}+$<br>$\int_0^t\{(a_{40}^{'})^{(8)}\left\|G_{40}^{(1)}-G_{40}^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}+$<br>$(a_{40}^{''})^{(8)}\left(T_{41}^{(1)},s_{(40)}\right)\left\|G_{40}^{(1)}-G_{40}^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{(\widehat{M}_{40})^{(8)}s_{(40)}}+$<br>$G_{40}^{(2)}\left\\|(a_{40}^{''})^{(8)}\left(T_{41}^{(1)},s_{(40)}\right)-(a_{40}^{''})^{(8)}\left(T_{41}^{(2)},s_{(40)}\right)\right\\|\ e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{(\widehat{M}_{40})^{(8)}s_{(40)}}\}ds_{(40)}$ | 271 |
| Where $s_{(40)}$ represents integrand that is integrated over the interval $[0,t]$<br><br>From the hypotheses it follows | 272 |
| $\left\|(G_{43})^{(1)}-(G_{43})^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}t}\le$<br>$\frac{1}{(\widehat{M}_{40})^{(8)}}\left((a_{40})^{(8)}+(a_{40}^{'})^{(8)}+(\widehat{A}_{40})^{(8)}+(\widehat{P}_{40})^{(8)}(\widehat{k}_{40})^{(8)}\right)d\left(((G_{43})^{(1)},(T_{43})^{(1)};\ (G_{43})^{(2)},(T_{43})^{(2)})\right)$<br><br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 273 |
| **Remark 36:** The fact that we supposed $(a_{40}^{''})^{(8)}$ and $(b_{40}^{''})^{(8)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{40})^{(8)}e^{(\widehat{M}_{40})^{(8)}t}$ *and* $(\widehat{Q}_{40})^{(8)}e^{(\widehat{M}_{40})^{(8)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(8)}$ and $(b_i^{''})^{(8)}, i=40,41,42$ depend only on $T_{41}$ and respectively on $(G_{43})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 274 |
| **Remark 37** There does not exist any $t$ where $G_i(t)=0$ *and* $T_i(t)=0$<br><br>it results<br><br>$G_i(t)\ge G_i^0e^{\left[-\int_0^t\{(a_i^{'})^{(8)}-(a_i^{''})^{(8)}(T_{41}(s_{(40)}),s_{(40)})\}ds_{(40)}\right]}\ge 0$<br><br>$T_i(t)\ge T_i^0e^{\left(-(b_i^{'})^{(8)}t\right)}>0$ for t > 0 | 275 |
| **Definition of** $\left((\widehat{M}_{40})^{(8)}\right)_1,\left((\widehat{M}_{40})^{(8)}\right)_2$ *and* $\left((\widehat{M}_{40})^{(8)}\right)_3$ : | 276 |

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| **Remark 38:** if $G_{40}$ is bounded, the same property have also $G_{41}\ and\ G_{42}$ . indeed if<br><br>$G_{40} < (\widehat{M}_{40})^{(8)}$ it follows $\frac{dG_{41}}{dt} \leq \left((\widehat{M}_{40})^{(8)}\right)_1 - (a'_{41})^{(8)}G_{41}$ and by integrating<br><br>$G_{41} \leq \left((\widehat{M}_{40})^{(8)}\right)_2 = G_{41}^0 + 2(a_{41})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_1/(a'_{41})^{(8)}$<br><br>In the same way , one can obtain<br><br>$G_{42} \leq \left((\widehat{M}_{40})^{(8)}\right)_3 = G_{42}^0 + 2(a_{42})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_2/(a'_{42})^{(8)}$<br><br>If $G_{41}\ or\ G_{42}$ is bounded, the same property follows for $G_{40}$ , $G_{42}$ and $G_{40}$ , $G_{41}$ respectively. | |
| **Remark 39:** If $G_{40}$ *is* bounded, from below, the same property holds for $G_{41}\ and\ G_{42}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{41}$ is bounded from below. | 277 |
| **Remark 40:** If $T_{40}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(8)}((G_{43})(t),t)) = (b'_{41})^{(8)}$ then $T_{41} \to \infty$.<br><br>**Definition of** $(m)^{(8)}$ and $\varepsilon_8$ :<br><br>Indeed let $t_8$ be so that for $t > t_8$<br><br>$(b_{41})^{(8)} - (b_i^{''})^{(8)}\left((G_{43})(t),t\right) < \varepsilon_8, T_{40}(t) > (m)^{(8)}$ | 278 |
| Then $\frac{dT_{41}}{dt} \geq (a_{41})^{(8)}(m)^{(8)} - \varepsilon_8 T_{41}$ which leads to<br><br>$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{\varepsilon_8}\right)(1 - e^{-\varepsilon_8 t}) + T_{41}^0 e^{-\varepsilon_8 t}$ If we take t such that $e^{-\varepsilon_8 t} = \frac{1}{2}$ it results<br><br>$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{2}\right),\ t = log\frac{2}{\varepsilon_8}$ By taking now $\varepsilon_8$ sufficiently small one sees that $T_{41}$ is unbounded.<br>The same property holds for $T_{42}$ if $\lim_{t\to\infty}(b_{42}^{''})^{(8)}\left((G_{43})(t),t(t),t\right) = (b'_{42})^{(8)}$ | 279 |
| It is now sufficient to take $\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}$ , $\frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}} < 1$ and to choose $(\widehat{P}_{44})^{(9)}$ and $(\widehat{Q}_{44})^{(9)}$ large to have | 279 A |
| $\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[(\widehat{P}_{44})^{(9)} + \left((\widehat{P}_{44})^{(9)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{44})^{(9)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{44})^{(9)}$ | |
| $\frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[\left((\widehat{Q}_{44})^{(9)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{44})^{(9)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{44})^{(9)}\right] \leq (\widehat{Q}_{44})^{(9)}$ | |

In order that the operator $\mathcal{A}^{(9)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying 39,35,36 into itself

The operator $\mathcal{A}^{(9)}$ is a contraction with respect to the metric

$$d\left(((G_{47})^{(1)},(T_{47})^{(1)}),((G_{47})^{(2)},(T_{47})^{(2)})\right)=$$
$$sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t)-G_i^{(2)}(t)\right|e^{-(\widehat{M}_{44})^{(9)}t},\max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t)-T_i^{(2)}(t)\right|e^{-(\widehat{M}_{44})^{(9)}t}\}$$

Indeed if we denote

**Definition of** $\widetilde{(G_{47})},\widetilde{(T_{47})}:\left(\widetilde{(G_{47})},\widetilde{(T_{47})}\right)=\mathcal{A}^{(9)}\left((G_{47}),(T_{47})\right)$
It results

$$\left|\tilde{G}_{44}^{(1)}-\tilde{G}_i^{(2)}\right|\le\int_0^t(a_{44})^{(9)}\left|G_{45}^{(1)}-G_{45}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\,ds_{(44)}+$$
$$\int_0^t\{(a_{44}^{'})^{(9)}\left|G_{44}^{(1)}-G_{44}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}+$$
$$(a_{44}^{''})^{(9)}\left(T_{45}^{(1)},s_{(44)}\right)\left|G_{44}^{(1)}-G_{44}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}+$$
$$G_{44}^{(2)}\left|(a_{44}^{''})^{(9)}\left(T_{45}^{(1)},s_{(44)}\right)-(a_{44}^{''})^{(9)}\left(T_{45}^{(2)},s_{(44)}\right)\right|\ e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\}ds_{(44)}$$

Where $s_{(44)}$ represents integrand that is integrated over the interval $[0,\mathrm{t}]$

From the hypotheses on 45,46,47,28 and 29 it follows

$$\left|(G_{47})^{(1)}-G^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}t}\le$$
$$\frac{1}{(\widehat{M}_{44})^{(9)}}\left((a_{44})^{(9)}+(a_{44}^{'})^{(9)}+(\widehat{A}_{44})^{(9)}+(\widehat{P}_{44})^{(9)}(\widehat{k}_{44})^{(9)}\right)d\left(((G_{47})^{(1)},(T_{47})^{(1)};\ (G_{47})^{(2)},(T_{47})^{(2)})\right)$$

And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis (39,35,36) the result follows

**Remark 41:** The fact that we supposed $(a_{44}^{''})^{(9)}$ and $(b_{44}^{''})^{(9)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{44})^{(9)}e^{(\widehat{M}_{44})^{(9)}t}$ *and* $(\widehat{Q}_{44})^{(9)}e^{(\widehat{M}_{44})^{(9)}t}$ respectively of $\mathbb{R}_+$.

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(9)}$ and $(b_i^{''})^{(9)}, i=44{,}45{,}46$ depend only on $\mathrm{T}_{45}$ and respectively on $(G_{47})$(*and not on* $t$) and hypothesis can replaced by a usual Lipschitz condition.

**Remark 42:** There does not exist any $t$ where $G_i(t)=0$ *and* $T_i(t)=0$

From 99 to 44 it results

$$G_i(t)\ge G_i^0e^{\left[-\int_0^t\{(a_i^{'})^{(9)}-(a_i^{''})^{(9)}(T_{45}(s_{(44)}),s_{(44)})\}ds_{(44)}\right]}\ge 0$$

$$T_i(t)\ge T_i^0e^{(-(b_i^{'})^{(9)}t)}>0\quad\text{for t}>0$$

**Definition of** $\left((\widehat{M}_{44})^{(9)}\right)_1,\left((\widehat{M}_{44})^{(9)}\right)_2$ *and* $\left((\widehat{M}_{44})^{(9)}\right)_3$ :

**Remark 43:** if $G_{44}$ is bounded, the same property have also $G_{45}\ and\ G_{46}$ . indeed if

$G_{44} < (\widehat{M}_{44})^{(9)}$ it follows $\frac{dG_{45}}{dt} \leq \left((\widehat{M}_{44})^{(9)}\right)_1 - (a_{45}^{'})^{(9)} G_{45}$ and by integrating

$G_{45} \leq \left((\widehat{M}_{44})^{(9)}\right)_2 = G_{45}^0 + 2(a_{45})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_1/(a_{45}^{'})^{(9)}$

In the same way , one can obtain

$G_{46} \leq \left((\widehat{M}_{44})^{(9)}\right)_3 = G_{46}^0 + 2(a_{46})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_2/(a_{46}^{'})^{(9)}$

If $G_{45}\ or\ G_{46}$ is bounded, the same property follows for $G_{44}$ , $G_{46}$ and $G_{44}$ , $G_{45}$ respectively.

**Remark 44:** If $G_{44}$ *is* bounded, from below, the same property holds for $G_{45}\ and\ G_{46}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{45}$ is bounded from below.

**Remark 45:** If $T_{44}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(9)}\left((G_{47})(t),t\right) = (b_{45}^{'})^{(9)}$ then $T_{45} \to \infty$.

**Definition of** $(m)^{(9)}$ and $\varepsilon_9$ :

Indeed let $t_9$ be so that for $t > t_9$

$$(b_{45})^{(9)} - (b_i^{''})^{(9)}\left((G_{47})(t),t\right) < \varepsilon_9, T_{44}(t) > (m)^{(9)}$$

Then $\frac{dT_{45}}{dt} \geq (a_{45})^{(9)}(m)^{(9)} - \varepsilon_9 T_{45}$ which leads to

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{\varepsilon_9}\right)(1 - e^{-\varepsilon_9 t}) + T_{45}^0 e^{-\varepsilon_9 t}$ If we take t such that $e^{-\varepsilon_9 t} = \frac{1}{2}$ it results

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{2}\right),\quad t = log\frac{2}{\varepsilon_9}$ By taking now $\varepsilon_9$ sufficiently small one sees that $T_{45}$ is unbounded.
The same property holds for $T_{46}$ if $\lim_{t\to\infty}(b_{46}^{''})^{(9)}\left((G_{47})(t),t\right) = (b_{46}^{'})^{(9)}$
We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 92

**Behavior of the solutions of equation** 280

**Theorem** If we denote and define

**Definition of** $(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ :

$(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ four constants satisfying
$-(\sigma_2)^{(1)} \leq -(a_{13}^{'})^{(1)} + (a_{14}^{'})^{(1)} - (a_{13}^{''})^{(1)}(T_{14},t) + (a_{14}^{''})^{(1)}(T_{14},t) \leq -(\sigma_1)^{(1)}$

$-(\tau_2)^{(1)} \leq -(b_{13}^{'})^{(1)} + (b_{14}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G,t) - (b_{14}^{''})^{(1)}(G,t) \leq -(\tau_1)^{(1)}$

**Definition of** $(v_1)^{(1)}, (v_2)^{(1)}, (u_1)^{(1)}, (u_2)^{(1)}, v^{(1)}, u^{(1)}$ : 281

By $(v_1)^{(1)} > 0$ , $(v_2)^{(1)} < 0$ and respectively $(u_1)^{(1)} > 0$ , $(u_2)^{(1)} < 0$ the roots of the equations
$(a_{14})^{(1)}\left(v^{(1)}\right)^2 + (\sigma_1)^{(1)}v^{(1)} - (a_{13})^{(1)} = 0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2 + (\tau_1)^{(1)}u^{(1)} - (b_{13})^{(1)} = 0$

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| **Definition of** $(\bar{v}_1)^{(1)},,(\bar{v}_2)^{(1)},(\bar{u}_1)^{(1)},(\bar{u}_2)^{(1)}$ :<br><br>By $(\bar{v}_1)^{(1)}>0\,,(\bar{v}_2)^{(1)}<0$ and respectively $(\bar{u}_1)^{(1)}>0\,,(\bar{u}_2)^{(1)}<0$ the roots of the equations $(a_{14})^{(1)}\left(v^{(1)}\right)^2+(\sigma_2)^{(1)}v^{(1)}-(a_{13})^{(1)}=0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2+(\tau_2)^{(1)}u^{(1)}-(b_{13})^{(1)}=0$ | 282 |
| **Definition of** $(m_1)^{(1)}\,,(m_2)^{(1)}\,,(\mu_1)^{(1)},(\mu_2)^{(1)},(\nu_0)^{(1)}$ :-<br><br>If we define $(m_1)^{(1)}\,,(m_2)^{(1)}\,,(\mu_1)^{(1)},(\mu_2)^{(1)}$ by<br>$(m_2)^{(1)}=(\nu_0)^{(1)},(m_1)^{(1)}=(\nu_1)^{(1)},\ if\ (\nu_0)^{(1)}<(\nu_1)^{(1)}$<br><br>$(m_2)^{(1)}=(\nu_1)^{(1)},(m_1)^{(1)}=(\bar{\nu}_1)^{(1)}\,,if\ (\nu_1)^{(1)}<(\nu_0)^{(1)}<(\bar{\nu}_1)^{(1)},$<br><br>and $\boxed{(\nu_0)^{(1)}=\frac{G_{13}^0}{G_{14}^0}}$<br><br>$(\,m_2)^{(1)}=(\nu_1)^{(1)},(m_1)^{(1)}=(\nu_0)^{(1)},\ if\ (\bar{\nu}_1)^{(1)}<(\nu_0)^{(1)}$ | 283 |
| and analogously<br><br>$(\mu_2)^{(1)}=(u_0)^{(1)},(\mu_1)^{(1)}=(u_1)^{(1)},\ if\ (u_0)^{(1)}<(u_1)^{(1)}$<br><br>$(\mu_2)^{(1)}=(u_1)^{(1)},(\mu_1)^{(1)}=(\bar{u}_1)^{(1)}\,,if\ (u_1)^{(1)}<(u_0)^{(1)}<(\bar{u}_1)^{(1)},$<br><br>and $\boxed{(u_0)^{(1)}=\frac{T_{13}^0}{T_{14}^0}}$<br><br>$(\,\mu_2)^{(1)}=(u_1)^{(1)},(\mu_1)^{(1)}=(u_0)^{(1)},if\ (\bar{u}_1)^{(1)}<(u_0)^{(1)}$ where $(u_1)^{(1)},(\bar{u}_1)^{(1)}$<br><br>are defined | 284 |
| Then the solution of global equations satisfies the inequalities<br><br>$G_{13}^0e^{((S_1)^{(1)}-(p_{13})^{(1)})t}\le G_{13}(t)\le G_{13}^0e^{(S_1)^{(1)}t}$<br><br>where $(p_i)^{(1)}$ is defined by equation<br><br>$\frac{1}{(m_1)^{(1)}}G_{13}^0e^{((S_1)^{(1)}-(p_{13})^{(1)})t}\le G_{14}(t)\le\frac{1}{(m_2)^{(1)}}G_{13}^0e^{(S_1)^{(1)}t}$ | 285 |
| $(\frac{(a_{15})^{(1)}G_{13}^0}{(m_1)^{(1)}((S_1)^{(1)}-(p_{13})^{(1)}-(S_2)^{(1)})}\left[e^{((S_1)^{(1)}-(p_{13})^{(1)})t}-e^{-(S_2)^{(1)}t}\right]+G_{15}^0e^{-(S_2)^{(1)}t}\le G_{15}(t)\le$<br><br>$\frac{(a_{15})^{(1)}G_{13}^0}{(m_2)^{(1)}((S_1)^{(1)}-(a_{15}^{'})^{(1)})}[e^{(S_1)^{(1)}t}-e^{-(a_{15}^{'})^{(1)}t}]+\ G_{15}^0e^{-(a_{15}^{'})^{(1)}t})$ | 286 |
| $\boxed{T_{13}^0e^{(R_1)^{(1)}t}\le T_{13}(t)\le T_{13}^0e^{((R_1)^{(1)}+(r_{13})^{(1)})t}}$ | 287 |
| $\frac{1}{(\mu_1)^{(1)}}T_{13}^0e^{(R_1)^{(1)}t}\le T_{13}(t)\le\frac{1}{(\mu_2)^{(1)}}T_{13}^0e^{((R_1)^{(1)}+(r_{13})^{(1)})t}$ | 288 |

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| $\frac{(b_{15})^{(1)}T_{13}^{0}}{(\mu_1)^{(1)}((R_1)^{(1)}-(b_{15}^{'})^{(1)})}\left[e^{(R_1)^{(1)}t}-e^{-(b_{15}^{'})^{(1)}t}\right]+T_{15}^{0}e^{-(b_{15}^{'})^{(1)}t}\leq T_{15}(t)\leq$ $\frac{(a_{15})^{(1)}T_{13}^{0}}{(\mu_2)^{(1)}((R_1)^{(1)}+(r_{13})^{(1)}+(R_2)^{(1)})}\left[e^{((R_1)^{(1)}+(r_{13})^{(1)})t}-e^{-(R_2)^{(1)}t}\right]+T_{15}^{0}e^{-(R_2)^{(1)}t}$ | 289 |
| **Definition of** $(S_1)^{(1)},(S_2)^{(1)},(R_1)^{(1)},(R_2)^{(1)}$:- Where $(S_1)^{(1)}=(a_{13})^{(1)}(m_2)^{(1)}-(a_{13}^{'})^{(1)}$ $(S_2)^{(1)}=(a_{15})^{(1)}-(p_{15})^{(1)}$ $(R_1)^{(1)}=(b_{13})^{(1)}(\mu_2)^{(1)}-(b_{13}^{'})^{(1)}$ $(R_2)^{(1)}=(b_{15}^{'})^{(1)}-(r_{15})^{(1)}$ | 290 |
| **Behavior of the solutions of equation** **Theorem 2:** If we denote and define | 291 |
| **Definition of** $(\sigma_1)^{(2)},(\sigma_2)^{(2)},(\tau_1)^{(2)},(\tau_2)^{(2)}$ : $(\sigma_1)^{(2)},(\sigma_2)^{(2)},(\tau_1)^{(2)},(\tau_2)^{(2)}$ four constants satisfying | 292 |
| $-(\sigma_2)^{(2)}\leq-(a_{16}^{'})^{(2)}+(a_{17}^{'})^{(2)}-(a_{16}^{''})^{(2)}(\mathrm{T}_{17},t)+(a_{17}^{''})^{(2)}(\mathrm{T}_{17},t)\leq-(\sigma_1)^{(2)}$ | 293 |
| $-(\tau_2)^{(2)}\leq-(b_{16}^{'})^{(2)}+(b_{17}^{'})^{(2)}-(b_{16}^{''})^{(2)}\big((G_{19}),t\big)-(b_{17}^{''})^{(2)}\big((G_{19}),t\big)\leq-(\tau_1)^{(2)}$ | 294 |
| **Definition of** $(\nu_1)^{(2)},(\nu_2)^{(2)},(u_1)^{(2)},(u_2)^{(2)}$ : | 295 |
| By $(\nu_1)^{(2)}>0$ , $(\nu_2)^{(2)}<0$ and respectively $(u_1)^{(2)}>0$ , $(u_2)^{(2)}<0$ the roots | 296 |
| of the equations $(a_{17})^{(2)}\big(\nu^{(2)}\big)^2+(\sigma_1)^{(2)}\nu^{(2)}-(a_{16})^{(2)}=0$ | 297 |
| and $(b_{14})^{(2)}\big(u^{(2)}\big)^2+(\tau_1)^{(2)}u^{(2)}-(b_{16})^{(2)}=0$ and | 298 |
| **Definition of** $(\bar{\nu}_1)^{(2)},,(\bar{\nu}_2)^{(2)},(\bar{u}_1)^{(2)},(\bar{u}_2)^{(2)}$ : | 299 |
| By $(\bar{\nu}_1)^{(2)}>0$ , $(\bar{\nu}_2)^{(2)}<0$ and respectively $(\bar{u}_1)^{(2)}>0$ , $(\bar{u}_2)^{(2)}<0$ the | 300 |
| roots of the equations $(a_{17})^{(2)}\big(\nu^{(2)}\big)^2+(\sigma_2)^{(2)}\nu^{(2)}-(a_{16})^{(2)}=0$ | 301 |
| and $(b_{17})^{(2)}\big(u^{(2)}\big)^2+(\tau_2)^{(2)}u^{(2)}-(b_{16})^{(2)}=0$ | 302 |
| **Definition of** $(m_1)^{(2)},(m_2)^{(2)},(\mu_1)^{(2)},(\mu_2)^{(2)}$ :- | 303 |
| If we define $(m_1)^{(2)},(m_2)^{(2)},(\mu_1)^{(2)},(\mu_2)^{(2)}$ by | 304 |
| $(m_2)^{(2)}=(\nu_0)^{(2)},(m_1)^{(2)}=(\nu_1)^{(2)}$, ***if*** $(\nu_0)^{(2)}<(\nu_1)^{(2)}$ | 305 |

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| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ , ***if*** $(\nu_1)^{(2)} < (\nu_0)^{(2)} < (\bar{\nu}_1)^{(2)}$, and $\boxed{(\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0}}$ | 306 |
| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\nu_0)^{(2)}$, ***if*** $(\bar{\nu}_1)^{(2)} < (\nu_0)^{(2)}$ | 307 |
| and analogously<br>$(\mu_2)^{(2)} = (u_0)^{(2)}, (\mu_1)^{(2)} = (u_1)^{(2)}$, ***if*** $(u_0)^{(2)} < (u_1)^{(2)}$<br>$(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (\bar{u}_1)^{(2)}$ , ***if*** $(u_1)^{(2)} < (u_0)^{(2)} < (\bar{u}_1)^{(2)}$, and $\boxed{(u_0)^{(2)} = \frac{T_{16}^0}{T_{17}^0}}$ | 308 |
| $(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (u_0)^{(2)}$, ***if*** $(\bar{u}_1)^{(2)} < (u_0)^{(2)}$ | 309 |
| Then the solution of global equations satisfies the inequalities<br>$G_{16}^0 e^{((S_1)^{(2)} - (p_{16})^{(2)})t} \le G_{16}(t) \le G_{16}^0 e^{(S_1)^{(2)} t}$ | 310 |
| $(p_i)^{(2)}$ is defined by equation | |
| $\frac{1}{(m_1)^{(2)}} G_{16}^0 e^{((S_1)^{(2)} - (p_{16})^{(2)})t} \le G_{17}(t) \le \frac{1}{(m_2)^{(2)}} G_{16}^0 e^{(S_1)^{(2)} t}$ | 311 |
| $\left(\frac{(a_{18})^{(2)} G_{16}^0}{(m_1)^{(2)}((S_1)^{(2)} - (p_{16})^{(2)} - (S_2)^{(2)})}\left[e^{((S_1)^{(2)} - (p_{16})^{(2)})t} - e^{-(S_2)^{(2)} t}\right] + G_{18}^0 e^{-(S_2)^{(2)} t} \le G_{18}(t) \le \right.$<br>$\left. \frac{(a_{18})^{(2)} G_{16}^0}{(m_2)^{(2)}((S_1)^{(2)} - (a'_{18})^{(2)})}\left[e^{(S_1)^{(2)} t} - e^{-(a'_{18})^{(2)} t}\right] + G_{18}^0 e^{-(a'_{18})^{(2)} t}\right)$ | 312 |
| $\boxed{T_{16}^0 e^{(R_1)^{(2)} t} \le T_{16}(t) \le T_{16}^0 e^{((R_1)^{(2)} + (r_{16})^{(2)})t}}$ | 313 |
| $\frac{1}{(\mu_1)^{(2)}} T_{16}^0 e^{(R_1)^{(2)} t} \le T_{16}(t) \le \frac{1}{(\mu_2)^{(2)}} T_{16}^0 e^{((R_1)^{(2)} + (r_{16})^{(2)})t}$ | 314 |
| $\frac{(b_{18})^{(2)} T_{16}^0}{(\mu_1)^{(2)}((R_1)^{(2)} - (b'_{18})^{(2)})}\left[e^{(R_1)^{(2)} t} - e^{-(b'_{18})^{(2)} t}\right] + T_{18}^0 e^{-(b'_{18})^{(2)} t} \le T_{18}(t) \le$<br>$\frac{(a_{18})^{(2)} T_{16}^0}{(\mu_2)^{(2)}((R_1)^{(2)} + (r_{16})^{(2)} + (R_2)^{(2)})}\left[e^{((R_1)^{(2)} + (r_{16})^{(2)})t} - e^{-(R_2)^{(2)} t}\right] + T_{18}^0 e^{-(R_2)^{(2)} t}$ | 315 |
| **Definition of** $(S_1)^{(2)}, (S_2)^{(2)}, (R_1)^{(2)}, (R_2)^{(2)}$:- | 316 |
| Where $(S_1)^{(2)} = (a_{16})^{(2)}(m_2)^{(2)} - (a'_{16})^{(2)}$<br>$(S_2)^{(2)} = (a_{18})^{(2)} - (p_{18})^{(2)}$ | 317 |
| $(R_1)^{(2)} = (b_{16})^{(2)}(\mu_2)^{(1)} - (b'_{16})^{(2)}$ | 318 |

$(R_2)^{(2)} = (b_{18}')^{(2)} - (r_{18})^{(2)}$

**Behavior of the solutions** 319

**Theorem 3:** If we denote and define

**Definition of** $(\sigma_1)^{(3)}, (\sigma_2)^{(3)}, (\tau_1)^{(3)}, (\tau_2)^{(3)}$ :

$(\sigma_1)^{(3)}, (\sigma_2)^{(3)}, (\tau_1)^{(3)}, (\tau_2)^{(3)}$ four constants satisfying

$$-(\sigma_2)^{(3)} \le -(a_{20}')^{(3)} + (a_{21}')^{(3)} - (a_{20}'')^{(3)}(T_{21}, t) + (a_{21}'')^{(3)}(T_{21}, t) \le -(\sigma_1)^{(3)}$$

$$-(\tau_2)^{(3)} \le -(b_{20}')^{(3)} + (b_{21}')^{(3)} - (b_{20}'')^{(3)}(G_{23}, t) - (b_{21}'')^{(3)}((G_{23}), t) \le -(\tau_1)^{(3)}$$

**Definition of** $(v_1)^{(3)}, (v_2)^{(3)}, (u_1)^{(3)}, (u_2)^{(3)}$ : 320

By $(v_1)^{(3)} > 0$ , $(v_2)^{(3)} < 0$ and respectively $(u_1)^{(3)} > 0$ , $(u_2)^{(3)} < 0$ the roots of the equations $(a_{21})^{(3)}\left(v^{(3)}\right)^2 + (\sigma_1)^{(3)}v^{(3)} - (a_{20})^{(3)} = 0$

and $(b_{21})^{(3)}\left(u^{(3)}\right)^2 + (\tau_1)^{(3)}u^{(3)} - (b_{20})^{(3)} = 0$ and

By $(\bar{v}_1)^{(3)} > 0$ , $(\bar{v}_2)^{(3)} < 0$ and respectively $(\bar{u}_1)^{(3)} > 0$ , $(\bar{u}_2)^{(3)} < 0$ the

roots of the equations $(a_{21})^{(3)}\left(v^{(3)}\right)^2 + (\sigma_2)^{(3)}v^{(3)} - (a_{20})^{(3)} = 0$

and $(b_{21})^{(3)}\left(u^{(3)}\right)^2 + (\tau_2)^{(3)}u^{(3)} - (b_{20})^{(3)} = 0$

**Definition of** $(m_1)^{(3)}, (m_2)^{(3)}, (\mu_1)^{(3)}, (\mu_2)^{(3)}$ :- 321

If we define $(m_1)^{(3)}, (m_2)^{(3)}, (\mu_1)^{(3)}, (\mu_2)^{(3)}$ by

$(m_2)^{(3)} = (v_0)^{(3)}, (m_1)^{(3)} = (v_1)^{(3)}$, ***if*** $(v_0)^{(3)} < (v_1)^{(3)}$

$(m_2)^{(3)} = (v_1)^{(3)}, (m_1)^{(3)} = (\bar{v}_1)^{(3)}$ , ***if*** $(v_1)^{(3)} < (v_0)^{(3)} < (\bar{v}_1)^{(3)}$,

and $\boxed{(v_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0}}$

$(m_2)^{(3)} = (v_1)^{(3)}, (m_1)^{(3)} = (v_0)^{(3)}$, ***if*** $(\bar{v}_1)^{(3)} < (v_0)^{(3)}$

and analogously 322

$(\mu_2)^{(3)} = (u_0)^{(3)}, (\mu_1)^{(3)} = (u_1)^{(3)}$, ***if*** $(u_0)^{(3)} < (u_1)^{(3)}$

$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (\bar{u}_1)^{(3)}$ , ***if*** $(u_1)^{(3)} < (u_0)^{(3)} < (\bar{u}_1)^{(3)}$, and $\boxed{(u_0)^{(3)} = \frac{T_{20}^0}{T_{21}^0}}$

$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (u_0)^{(3)}$, ***if*** $(\bar{u}_1)^{(3)} < (u_0)^{(3)}$

Then the solution of global equations satisfies the inequalities

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| $G_{20}^{0}e^{((S_1)^{(3)}-(p_{20})^{(3)})t} \le G_{20}(t) \le G_{20}^{0}e^{(S_1)^{(3)}t}$<br>$(p_i)^{(3)}$ is defined by equation | |
| $\frac{1}{(m_1)^{(3)}}G_{20}^{0}e^{((S_1)^{(3)}-(p_{20})^{(3)})t} \le G_{21}(t) \le \frac{1}{(m_2)^{(3)}}G_{20}^{0}e^{(S_1)^{(3)}t}$ | 323 |
| $(\frac{(a_{22})^{(3)}G_{20}^{0}}{(m_1)^{(3)}((S_1)^{(3)}-(p_{20})^{(3)}-(S_2)^{(3)})}\left[e^{((S_1)^{(3)}-(p_{20})^{(3)})t}-e^{-(S_2)^{(3)}t}\right]+G_{22}^{0}e^{-(S_2)^{(3)}t} \le G_{22}(t) \le$<br>$\frac{(a_{22})^{(3)}G_{20}^{0}}{(m_2)^{(3)}((S_1)^{(3)}-(a_{22}')^{(3)})}[e^{(S_1)^{(3)}t}-e^{-(a_{22}')^{(3)}t}]+G_{22}^{0}e^{-(a_{22}')^{(3)}t})$ | 324 |
| $\boxed{T_{20}^{0}e^{(R_1)^{(3)}t} \le T_{20}(t) \le T_{20}^{0}e^{((R_1)^{(3)}+(r_{20})^{(3)})t}}$ | 325 |
| $\frac{1}{(\mu_1)^{(3)}}T_{20}^{0}e^{(R_1)^{(3)}t} \le T_{20}(t) \le \frac{1}{(\mu_2)^{(3)}}T_{20}^{0}e^{((R_1)^{(3)}+(r_{20})^{(3)})t}$ | 326 |
| $\frac{(b_{22})^{(3)}T_{20}^{0}}{(\mu_1)^{(3)}((R_1)^{(3)}-(b_{22}')^{(3)})}\left[e^{(R_1)^{(3)}t}-e^{-(b_{22}')^{(3)}t}\right]+T_{22}^{0}e^{-(b_{22}')^{(3)}t} \le T_{22}(t) \le$<br>$\frac{(a_{22})^{(3)}T_{20}^{0}}{(\mu_2)^{(3)}((R_1)^{(3)}+(r_{20})^{(3)}+(R_2)^{(3)})}\left[e^{((R_1)^{(3)}+(r_{20})^{(3)})t}-e^{-(R_2)^{(3)}t}\right]+T_{22}^{0}e^{-(R_2)^{(3)}t}$ | 327 |
| **Definition of** $(S_1)^{(3)},(S_2)^{(3)},(R_1)^{(3)},(R_2)^{(3)}$:-<br>Where $(S_1)^{(3)}=(a_{20})^{(3)}(m_2)^{(3)}-(a_{20}')^{(3)}$<br>$(S_2)^{(3)}=(a_{22})^{(3)}-(p_{22})^{(3)}$<br>$(R_1)^{(3)}=(b_{20})^{(3)}(\mu_2)^{(3)}-(b_{20}')^{(3)}$<br>$(R_2)^{(3)}=(b_{22}')^{(3)}-(r_{22})^{(3)}$ | 328 |
| **Behavior of the solutions of equation**<br>**Theorem:** If we denote and define<br>**Definition of** $(\sigma_1)^{(4)},(\sigma_2)^{(4)},(\tau_1)^{(4)},(\tau_2)^{(4)}$ :<br>$(\sigma_1)^{(4)},(\sigma_2)^{(4)},(\tau_1)^{(4)},(\tau_2)^{(4)}$ four constants satisfying<br>$-(\sigma_2)^{(4)} \le -(a_{24}')^{(4)}+(a_{25}')^{(4)}-(a_{24}'')^{(4)}(T_{25},t)+(a_{25}'')^{(4)}(T_{25},t) \le -(\sigma_1)^{(4)}$<br>$-(\tau_2)^{(4)} \le -(b_{24}')^{(4)}+(b_{25}')^{(4)}-(b_{24}'')^{(4)}((G_{27}),t)-(b_{25}'')^{(4)}((G_{27}),t) \le -(\tau_1)^{(4)}$ | |
| **Definition of** $(v_1)^{(4)},(v_2)^{(4)},(u_1)^{(4)},(u_2)^{(4)},v^{(4)},u^{(4)}$ :<br>By $(v_1)^{(4)}>0$, $(v_2)^{(4)}<0$ and respectively $(u_1)^{(4)}>0$, $(u_2)^{(4)}<0$ the roots of the equations $(a_{25})^{(4)}(v^{(4)})^2+(\sigma_1)^{(4)}v^{(4)}-(a_{24})^{(4)}=0$<br>and $(b_{25})^{(4)}(u^{(4)})^2+(\tau_1)^{(4)}u^{(4)}-(b_{24})^{(4)}=0$ and | 329 |

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| **Definition of** $(\bar{v}_1)^{(4)},,(\bar{v}_2)^{(4)},(\bar{u}_1)^{(4)},(\bar{u}_2)^{(4)}$ :<br><br>By $(\bar{v}_1)^{(4)}>0$ , $(\bar{v}_2)^{(4)}<0$ and respectively $(\bar{u}_1)^{(4)}>0$ , $(\bar{u}_2)^{(4)}<0$ the roots of the equations $(a_{25})^{(4)}\left(v^{(4)}\right)^2+(\sigma_2)^{(4)}v^{(4)}-(a_{24})^{(4)}=0$ and $(b_{25})^{(4)}\left(u^{(4)}\right)^2+(\tau_2)^{(4)}u^{(4)}-(b_{24})^{(4)}=0$<br>**Definition of** $(m_1)^{(4)}$ , $(m_2)^{(4)}$ , $(\mu_1)^{(4)},(\mu_2)^{(4)},(v_0)^{(4)}$ :-<br><br>If we define $(m_1)^{(4)}$ , $(m_2)^{(4)}$ , $(\mu_1)^{(4)},(\mu_2)^{(4)}$ by<br><br>$(m_2)^{(4)}=(v_0)^{(4)},(m_1)^{(4)}=(v_1)^{(4)}$, ***if*** $(v_0)^{(4)}<(v_1)^{(4)}$<br><br>$(m_2)^{(4)}=(v_1)^{(4)},(m_1)^{(4)}=(\bar{v}_1)^{(4)}$ , ***if*** $(v_4)^{(4)}<(v_0)^{(4)}<(\bar{v}_1)^{(4)}$,<br>and $\boxed{(v_0)^{(4)}=\frac{G_{24}^0}{G_{25}^0}}$<br><br>$(m_2)^{(4)}=(v_4)^{(4)},(m_1)^{(4)}=(v_0)^{(4)}$, ***if*** $(\bar{v}_4)^{(4)}<(v_0)^{(4)}$ | 330 |
| and analogously<br><br>$(\mu_2)^{(4)}=(u_0)^{(4)},(\mu_1)^{(4)}=(u_1)^{(4)}$, ***if*** $(u_0)^{(4)}<(u_1)^{(4)}$<br><br>$(\mu_2)^{(4)}=(u_1)^{(4)},(\mu_1)^{(4)}=(\bar{u}_1)^{(4)}$ , ***if*** $(u_1)^{(4)}<(u_0)^{(4)}<(\bar{u}_1)^{(4)}$,<br>and $\boxed{(u_0)^{(4)}=\frac{T_{24}^0}{T_{25}^0}}$<br><br>$(\mu_2)^{(4)}=(u_1)^{(4)},(\mu_1)^{(4)}=(u_0)^{(4)}$, ***if*** $(\bar{u}_1)^{(4)}<(u_0)^{(4)}$ where $(u_1)^{(4)},(\bar{u}_1)^{(4)}$ | 331 |
| Then the solution of global equations satisfies the inequalities<br><br>$G_{24}^0e^{((S_1)^{(4)}-(p_{24})^{(4)})t}\le G_{24}(t)\le G_{24}^0e^{(S_1)^{(4)}t}$<br><br>where $(p_i)^{(4)}$ is defined by equation | 332 |
| $\frac{1}{(m_1)^{(4)}}G_{24}^0e^{((S_1)^{(4)}-(p_{24})^{(4)})t}\le G_{25}(t)\le\frac{1}{(m_2)^{(4)}}G_{24}^0e^{(S_1)^{(4)}t}$ | 333 |
| $\left(\frac{(a_{26})^{(4)}G_{24}^0}{(m_1)^{(4)}((S_1)^{(4)}-(p_{24})^{(4)}-(S_2)^{(4)})}\left[e^{((S_1)^{(4)}-(p_{24})^{(4)})t}-e^{-(S_2)^{(4)}t}\right]+G_{26}^0e^{-(S_2)^{(4)}t}\le G_{26}(t)\le\right.$<br>*(a26)4G240(m2)4(S1)4−(a26′)4e(S1)4t−e−(a26′)4t+ G260e−(a26′)4t* | 334 |
| $\boxed{T_{24}^0e^{(R_1)^{(4)}t}\le T_{24}(t)\le T_{24}^0e^{((R_1)^{(4)}+(r_{24})^{(4)})t}}$ | |
| $\frac{1}{(\mu_1)^{(4)}}T_{24}^0e^{(R_1)^{(4)}t}\le T_{24}(t)\le\frac{1}{(\mu_2)^{(4)}}T_{24}^0e^{((R_1)^{(4)}+(r_{24})^{(4)})t}$ | 335 |
| $\frac{(b_{26})^{(4)}T_{24}^0}{(\mu_1)^{(4)}((R_1)^{(4)}-(b_{26}^{'})^{(4)})}\left[e^{(R_1)^{(4)}t}-e^{-(b_{26}^{'})^{(4)}t}\right]+T_{26}^0e^{-(b_{26}^{'})^{(4)}t}\le T_{26}(t)\le$<br><br>$\frac{(a_{26})^{(4)}T_{24}^0}{(\mu_2)^{(4)}((R_1)^{(4)}+(r_{24})^{(4)}+(R_2)^{(4)})}\left[e^{((R_1)^{(4)}+(r_{24})^{(4)})t}-e^{-(R_2)^{(4)}t}\right]+T_{26}^0e^{-(R_2)^{(4)}t}$ | 336 |

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| **Definition of** $(S_1)^{(4)}, (S_2)^{(4)}, (R_1)^{(4)}, (R_2)^{(4)}$:-<br>Where $(S_1)^{(4)} = (a_{24})^{(4)}(m_2)^{(4)} - (a'_{24})^{(4)}$<br>$(S_2)^{(4)} = (a_{26})^{(4)} - (p_{26})^{(4)}$<br>$(R_1)^{(4)} = (b_{24})^{(4)}(\mu_2)^{(4)} - (b'_{24})^{(4)}$<br>$(R_2)^{(4)} = (b'_{26})^{(4)} - (r_{26})^{(4)}$ | 337 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ :<br>$(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ four constants satisfying<br>$-(\sigma_2)^{(5)} \le -(a'_{28})^{(5)} + (a'_{29})^{(5)} - (a''_{28})^{(5)}(T_{29}, t) + (a''_{29})^{(5)}(T_{29}, t) \le -(\sigma_1)^{(5)}$<br>$-(\tau_2)^{(5)} \le -(b'_{28})^{(5)} + (b'_{29})^{(5)} - (b''_{28})^{(5)}((G_{31}), t) - (b''_{29})^{(5)}((G_{31}), t) \le -(\tau_1)^{(5)}$ | 338 |
| **Definition of** $(v_1)^{(5)}, (v_2)^{(5)}, (u_1)^{(5)}, (u_2)^{(5)}, v^{(5)}, u^{(5)}$ :<br>By $(v_1)^{(5)} > 0$, $(v_2)^{(5)} < 0$ and respectively $(u_1)^{(5)} > 0$, $(u_2)^{(5)} < 0$ the roots of the equations $(a_{29})^{(5)}\left(v^{(5)}\right)^2 + (\sigma_1)^{(5)}v^{(5)} - (a_{28})^{(5)} = 0$<br>and $(b_{29})^{(5)}\left(u^{(5)}\right)^2 + (\tau_1)^{(5)}u^{(5)} - (b_{28})^{(5)} = 0$ and | 339 |
| **Definition of** $(\bar{v}_1)^{(5)},, (\bar{v}_2)^{(5)}, (\bar{u}_1)^{(5)}, (\bar{u}_2)^{(5)}$ :<br>By $(\bar{v}_1)^{(5)} > 0$, $(\bar{v}_2)^{(5)} < 0$ and respectively $(\bar{u}_1)^{(5)} > 0$, $(\bar{u}_2)^{(5)} < 0$ the roots of the equations $(a_{29})^{(5)}\left(v^{(5)}\right)^2 + (\sigma_2)^{(5)}v^{(5)} - (a_{28})^{(5)} = 0$<br>and $(b_{29})^{(5)}\left(u^{(5)}\right)^2 + (\tau_2)^{(5)}u^{(5)} - (b_{28})^{(5)} = 0$<br>**Definition of** $(m_1)^{(5)}, (m_2)^{(5)}, (\mu_1)^{(5)}, (\mu_2)^{(5)}, (v_0)^{(5)}$ :-<br>If we define $(m_1)^{(5)}, (m_2)^{(5)}, (\mu_1)^{(5)}, (\mu_2)^{(5)}$ by<br>$(m_2)^{(5)} = (v_0)^{(5)}, (m_1)^{(5)} = (v_1)^{(5)}$, ***if*** $(v_0)^{(5)} < (v_1)^{(5)}$<br>$(m_2)^{(5)} = (v_1)^{(5)}, (m_1)^{(5)} = (\bar{v}_1)^{(5)}$, ***if*** $(v_1)^{(5)} < (v_0)^{(5)} < (\bar{v}_1)^{(5)}$,<br>and $\boxed{(v_0)^{(5)} = \frac{G^0_{28}}{G^0_{29}}}$<br>$(m_2)^{(5)} = (v_1)^{(5)}, (m_1)^{(5)} = (v_0)^{(5)}$, ***if*** $(\bar{v}_1)^{(5)} < (v_0)^{(5)}$ | 340 |
| and analogously<br>$(\mu_2)^{(5)} = (u_0)^{(5)}, (\mu_1)^{(5)} = (u_1)^{(5)}$, ***if*** $(u_0)^{(5)} < (u_1)^{(5)}$ | 341 |

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| $(\mu_2)^{(5)} = (u_1)^{(5)}, (\mu_1)^{(5)} = (\bar{u}_1)^{(5)}$ , $\boldsymbol{if}$ $(u_1)^{(5)} < (u_0)^{(5)} < (\bar{u}_1)^{(5)}$, and $\boxed{(u_0)^{(5)} = \frac{T_{28}^0}{T_{29}^0}}$ $(\mu_2)^{(5)} = (u_1)^{(5)}, (\mu_1)^{(5)} = (u_0)^{(5)}, \boldsymbol{if}$ $(\bar{u}_1)^{(5)} < (u_0)^{(5)}$ where $(u_1)^{(5)}, (\bar{u}_1)^{(5)}$ | |
| Then the solution of global equations satisfies the inequalities $G_{28}^0 e^{((S_1)^{(5)} - (p_{28})^{(5)})t} \le G_{28}(t) \le G_{28}^0 e^{(S_1)^{(5)} t}$ where $(p_i)^{(5)}$ is defined by equation | 342 |
| $\frac{1}{(m_5)^{(5)}} G_{28}^0 e^{((S_1)^{(5)} - (p_{28})^{(5)})t} \le G_{29}(t) \le \frac{1}{(m_2)^{(5)}} G_{28}^0 e^{(S_1)^{(5)} t}$ | 343 |
| $\left(\frac{(a_{30})^{(5)} G_{28}^0}{(m_1)^{(5)}((S_1)^{(5)} - (p_{28})^{(5)} - (S_2)^{(5)})}\left[e^{((S_1)^{(5)} - (p_{28})^{(5)})t} - e^{-(S_2)^{(5)} t}\right] + G_{30}^0 e^{-(S_2)^{(5)} t} \le G_{30}(t) \le$ *(a30)5G280(m2)5(S1)5−(a30′)5e(S1)5t−e−(a30′)5t+ G300e−(a30′)5t* | 344 |
| $\boxed{T_{28}^0 e^{(R_1)^{(5)} t} \le T_{28}(t) \le T_{28}^0 e^{((R_1)^{(5)} + (r_{28})^{(5)})t}}$ | 345 |
| $\frac{1}{(\mu_1)^{(5)}} T_{28}^0 e^{(R_1)^{(5)} t} \le T_{28}(t) \le \frac{1}{(\mu_2)^{(5)}} T_{28}^0 e^{((R_1)^{(5)} + (r_{28})^{(5)})t}$ | 346 |
| $\frac{(b_{30})^{(5)} T_{28}^0}{(\mu_1)^{(5)}((R_1)^{(5)} - (b_{30}^{'})^{(5)})}\left[e^{(R_1)^{(5)} t} - e^{-(b_{30}^{'})^{(5)} t}\right] + T_{30}^0 e^{-(b_{30}^{'})^{(5)} t} \le T_{30}(t) \le$ $\frac{(a_{30})^{(5)} T_{28}^0}{(\mu_2)^{(5)}((R_1)^{(5)} + (r_{28})^{(5)} + (R_2)^{(5)})}\left[e^{((R_1)^{(5)} + (r_{28})^{(5)})t} - e^{-(R_2)^{(5)} t}\right] + T_{30}^0 e^{-(R_2)^{(5)} t}$ | 347 |
| **Definition of** $(S_1)^{(5)}, (S_2)^{(5)}, (R_1)^{(5)}, (R_2)^{(5)}$:- Where $(S_1)^{(5)} = (a_{28})^{(5)}(m_2)^{(5)} - (a_{28}^{'})^{(5)}$ $(S_2)^{(5)} = (a_{30})^{(5)} - (p_{30})^{(5)}$ $(R_1)^{(5)} = (b_{28})^{(5)}(\mu_2)^{(5)} - (b_{28}^{'})^{(5)}$ $(R_2)^{(5)} = (b_{30}^{'})^{(5)} - (r_{30})^{(5)}$ | 348 |
| **Behavior of the solutions of equation** **Theorem 2:** If we denote and define **Definition of** $(\sigma_1)^{(6)}, (\sigma_2)^{(6)}, (\tau_1)^{(6)}, (\tau_2)^{(6)}$ : $(\sigma_1)^{(6)}, (\sigma_2)^{(6)}, (\tau_1)^{(6)}, (\tau_2)^{(6)}$ four constants satisfying $-(\sigma_2)^{(6)} \le -(a_{32}^{'})^{(6)} + (a_{33}^{'})^{(6)} - (a_{32}^{''})^{(6)}(T_{33}, t) + (a_{33}^{''})^{(6)}(T_{33}, t) \le -(\sigma_1)^{(6)}$ | 349 |

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| $-(\tau_2)^{(6)} \le -(b'_{32})^{(6)} + (b'_{33})^{(6)} - (b''_{32})^{(6)}\big((G_{35}),t\big) - (b''_{33})^{(6)}\big((G_{35}),t\big) \le -(\tau_1)^{(6)}$ | |
| **Definition of** $(\nu_1)^{(6)},(\nu_2)^{(6)},(u_1)^{(6)},(u_2)^{(6)},\nu^{(6)},u^{(6)}$ :<br><br>By $(\nu_1)^{(6)} > 0$ , $(\nu_2)^{(6)} < 0$ and respectively $(u_1)^{(6)} > 0$ , $(u_2)^{(6)} < 0$ the roots of the equations $(a_{33})^{(6)}\big(\nu^{(6)}\big)^2 + (\sigma_1)^{(6)}\nu^{(6)} - (a_{32})^{(6)} = 0$<br>and $(b_{33})^{(6)}\big(u^{(6)}\big)^2 + (\tau_1)^{(6)}u^{(6)} - (b_{32})^{(6)} = 0$ and | 350 |
| **Definition of** $(\bar{\nu}_1)^{(6)},,(\bar{\nu}_2)^{(6)},(\bar{u}_1)^{(6)},(\bar{u}_2)^{(6)}$ :<br><br>By $(\bar{\nu}_1)^{(6)} > 0$ , $(\bar{\nu}_2)^{(6)} < 0$ and respectively $(\bar{u}_1)^{(6)} > 0$ , $(\bar{u}_2)^{(6)} < 0$ the roots of the equations $(a_{33})^{(6)}\big(\nu^{(6)}\big)^2 + (\sigma_2)^{(6)}\nu^{(6)} - (a_{32})^{(6)} = 0$<br>and $(b_{33})^{(6)}\big(u^{(6)}\big)^2 + (\tau_2)^{(6)}u^{(6)} - (b_{32})^{(6)} = 0$<br>**Definition of** $(m_1)^{(6)}$ , $(m_2)^{(6)}$ , $(\mu_1)^{(6)},(\mu_2)^{(6)},(\nu_0)^{(6)}$ :-<br><br>If we define $(m_1)^{(6)}$ , $(m_2)^{(6)}$ , $(\mu_1)^{(6)},(\mu_2)^{(6)}$ by<br><br>$(m_2)^{(6)} = (\nu_0)^{(6)}, (m_1)^{(6)} = (\nu_1)^{(6)}$, ***if*** $(\nu_0)^{(6)} < (\nu_1)^{(6)}$<br><br>$(m_2)^{(6)} = (\nu_1)^{(6)}, (m_1)^{(6)} = (\bar{\nu}_6)^{(6)}$ , ***if*** $(\nu_1)^{(6)} < (\nu_0)^{(6)} < (\bar{\nu}_1)^{(6)}$,<br>and $\boxed{(\nu_0)^{(6)} = \frac{G^0_{32}}{G^0_{33}}}$<br><br>$(m_2)^{(6)} = (\nu_1)^{(6)}, (m_1)^{(6)} = (\nu_0)^{(6)}$, ***if*** $(\bar{\nu}_1)^{(6)} < (\nu_0)^{(6)}$ | 351 |
| and analogously<br><br>$(\mu_2)^{(6)} = (u_0)^{(6)}, (\mu_1)^{(6)} = (u_1)^{(6)}$, ***if*** $(u_0)^{(6)} < (u_1)^{(6)}$<br><br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (\bar{u}_1)^{(6)}$ , ***if*** $(u_1)^{(6)} < (u_0)^{(6)} < (\bar{u}_1)^{(6)}$,<br>and $\boxed{(u_0)^{(6)} = \frac{T^0_{32}}{T^0_{33}}}$<br><br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (u_0)^{(6)}$, ***if*** $(\bar{u}_1)^{(6)} < (u_0)^{(6)}$ where $(u_1)^{(6)}, (\bar{u}_1)^{(6)}$ | 352 |
| Then the solution of global equations satisfies the inequalities<br><br>$G^0_{32}e^{((S_1)^{(6)}-(p_{32})^{(6)})t} \le G_{32}(t) \le G^0_{32}e^{(S_1)^{(6)}t}$<br><br>where $(p_i)^{(6)}$ is defined by equation | 353 |
| $\frac{1}{(m_1)^{(6)}}G^0_{32}e^{((S_1)^{(6)}-(p_{32})^{(6)})t} \le G_{33}(t) \le \frac{1}{(m_2)^{(6)}}G^0_{32}e^{(S_1)^{(6)}t}$ | 354 |
| $\left(\frac{(a_{34})^{(6)}G^0_{32}}{(m_1)^{(6)}((S_1)^{(6)}-(p_{32})^{(6)}-(S_2)^{(6)})}\left[e^{((S_1)^{(6)}-(p_{32})^{(6)})t} - e^{-(S_2)^{(6)}t}\right] + G^0_{34}e^{-(S_2)^{(6)}t} \le G_{34}(t) \le$<br>*(a34)6G320(m2)6(S1)6−(a34′)6e(S1)6t−e−(a34′)6t+ G340e−(a34′)6t* | 355 |

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| $T_{32}^{0}e^{(R_1)^{(6)}t} \le T_{32}(t) \le T_{32}^{0}e^{((R_1)^{(6)}+(r_{32})^{(6)})t}$ | 356 |
| $\frac{1}{(\mu_1)^{(6)}}T_{32}^{0}e^{(R_1)^{(6)}t} \le T_{32}(t) \le \frac{1}{(\mu_2)^{(6)}}T_{32}^{0}e^{((R_1)^{(6)}+(r_{32})^{(6)})t}$ | 357 |
| $\frac{(b_{34})^{(6)}T_{32}^{0}}{(\mu_1)^{(6)}((R_1)^{(6)}-(b_{34}^{'})^{(6)})}\left[e^{(R_1)^{(6)}t}-e^{-(b_{34}^{'})^{(6)}t}\right]+T_{34}^{0}e^{-(b_{34}^{'})^{(6)}t} \le T_{34}(t) \le$ $\frac{(a_{34})^{(6)}T_{32}^{0}}{(\mu_2)^{(6)}((R_1)^{(6)}+(r_{32})^{(6)}+(R_2)^{(6)})}\left[e^{((R_1)^{(6)}+(r_{32})^{(6)})t}-e^{-(R_2)^{(6)}t}\right]+T_{34}^{0}e^{-(R_2)^{(6)}t}$ | 358 |
| **Definition of** $(S_1)^{(6)},(S_2)^{(6)},(R_1)^{(6)},(R_2)^{(6)}$:- Where $(S_1)^{(6)}=(a_{32})^{(6)}(m_2)^{(6)}-(a_{32}^{'})^{(6)}$ $(S_2)^{(6)}=(a_{34})^{(6)}-(p_{34})^{(6)}$ $(R_1)^{(6)}=(b_{32})^{(6)}(\mu_2)^{(6)}-(b_{32}^{'})^{(6)}$ $(R_2)^{(6)}=(b_{34}^{'})^{(6)}-(r_{34})^{(6)}$ | 359 |
| **Behavior of the solutions of equation** **Theorem 2:** If we denote and define **Definition of** $(\sigma_1)^{(7)},(\sigma_2)^{(7)},(\tau_1)^{(7)},(\tau_2)^{(7)}$: $(\sigma_1)^{(7)},(\sigma_2)^{(7)},(\tau_1)^{(7)},(\tau_2)^{(7)}$ four constants satisfying $-(\sigma_2)^{(7)} \le -(a_{36}^{'})^{(7)}+(a_{37}^{'})^{(7)}-(a_{36}^{''})^{(7)}(T_{37},t)+(a_{37}^{''})^{(7)}(T_{37},t) \le -(\sigma_1)^{(7)}$ $-(\tau_2)^{(7)} \le -(b_{36}^{'})^{(7)}+(b_{37}^{'})^{(7)}-(b_{36}^{''})^{(7)}\big((G_{39}),t\big)-(b_{37}^{''})^{(7)}\big((G_{39}),t\big) \le -(\tau_1)^{(7)}$ | |
| **Definition of** $(v_1)^{(7)},(v_2)^{(7)},(u_1)^{(7)},(u_2)^{(7)},v^{(7)},u^{(7)}$: By $(v_1)^{(7)}>0,(v_2)^{(7)}<0$ and respectively $(u_1)^{(7)}>0,(u_2)^{(7)}<0$ the roots of the equations $(a_{37})^{(7)}\big(v^{(7)}\big)^2+(\sigma_1)^{(7)}v^{(7)}-(a_{36})^{(7)}=0$ and $(b_{37})^{(7)}\big(u^{(7)}\big)^2+(\tau_1)^{(7)}u^{(7)}-(b_{36})^{(7)}=0$ and | 361 |
| **Definition of** $(\bar{v}_1)^{(7)},,(\bar{v}_2)^{(7)},(\bar{u}_1)^{(7)},(\bar{u}_2)^{(7)}$: By $(\bar{v}_1)^{(7)}>0,(\bar{v}_2)^{(7)}<0$ and respectively $(\bar{u}_1)^{(7)}>0,(\bar{u}_2)^{(7)}<0$ the roots of the equations $(a_{37})^{(7)}\big(v^{(7)}\big)^2+(\sigma_2)^{(7)}v^{(7)}-(a_{36})^{(7)}=0$ and $(b_{37})^{(7)}\big(u^{(7)}\big)^2+(\tau_2)^{(7)}u^{(7)}-(b_{36})^{(7)}=0$ **Definition of** $(m_1)^{(7)},(m_2)^{(7)},(\mu_1)^{(7)},(\mu_2)^{(7)},(v_0)^{(7)}$:- | 362 |

If we define $(m_1)^{(7)}, (m_2)^{(7)}, (\mu_1)^{(7)}, (\mu_2)^{(7)}$ by

$(m_2)^{(7)} = (\nu_0)^{(7)}, (m_1)^{(7)} = (\nu_1)^{(7)}$, $\boldsymbol{if}\ (\nu_0)^{(7)} < (\nu_1)^{(7)}$

$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\bar{\nu}_1)^{(7)}$, $\boldsymbol{if}\ (\nu_1)^{(7)} < (\nu_0)^{(7)} < (\bar{\nu}_1)^{(7)}$,

and $\boxed{(\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0}}$

$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\nu_0)^{(7)}$, $\boldsymbol{if}\ (\bar{\nu}_1)^{(7)} < (\nu_0)^{(7)}$

and analogously

$(\mu_2)^{(7)} = (u_0)^{(7)}, (\mu_1)^{(7)} = (u_1)^{(7)}$, $\boldsymbol{if}\ (u_0)^{(7)} < (u_1)^{(7)}$

$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (\bar{u}_1)^{(7)}$, $\boldsymbol{if}\ (u_1)^{(7)} < (u_0)^{(7)} < (\bar{u}_1)^{(7)}$,

and $\boxed{(u_0)^{(7)} = \frac{T_{36}^0}{T_{37}^0}}$

$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (u_0)^{(7)}, \boldsymbol{if}\ (\bar{u}_1)^{(7)} < (u_0)^{(7)}$ where $(u_1)^{(7)}, (\bar{u}_1)^{(7)}$

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Then the solution of global equations satisfies the inequalities

$$G_{36}^0 e^{((S_1)^{(7)} - (p_{36})^{(7)})t} \le G_{36}(t) \le G_{36}^0 e^{(S_1)^{(7)} t}$$

where $(p_i)^{(7)}$ is defined by equation

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$$\frac{1}{(m_7)^{(7)}} G_{36}^0 e^{((S_1)^{(7)} - (p_{36})^{(7)})t} \le G_{37}(t) \le \frac{1}{(m_2)^{(7)}} G_{36}^0 e^{(S_1)^{(7)} t}$$

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$$\left(\frac{(a_{38})^{(7)} G_{36}^0}{(m_1)^{(7)}((S_1)^{(7)} - (p_{36})^{(7)} - (S_2)^{(7)})}\left[e^{((S_1)^{(7)} - (p_{36})^{(7)})t} - e^{-(S_2)^{(7)} t}\right] + G_{38}^0 e^{-(S_2)^{(7)} t} \le G_{38}(t) \le \right.$$

$$\left. \frac{(a_{38})^{(7)} G_{36}^0}{(m_2)^{(7)}((S_1)^{(7)} - (a_{38}')^{(7)})}\left[e^{(S_1)^{(7)} t} - e^{-(a_{38}')^{(7)} t}\right] + G_{38}^0 e^{-(a_{38}')^{(7)} t}\right)$$

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$$\boxed{T_{36}^0 e^{(R_1)^{(7)} t} \le T_{36}(t) \le T_{36}^0 e^{((R_1)^{(7)} + (r_{36})^{(7)})t}}$$

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$$\frac{1}{(\mu_1)^{(7)}} T_{36}^0 e^{(R_1)^{(7)} t} \le T_{36}(t) \le \frac{1}{(\mu_2)^{(7)}} T_{36}^0 e^{((R_1)^{(7)} + (r_{36})^{(7)})t}$$

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$$\frac{(b_{38})^{(7)} T_{36}^0}{(\mu_1)^{(7)}((R_1)^{(7)} - (b_{38}')^{(7)})}\left[e^{(R_1)^{(7)} t} - e^{-(b_{38}')^{(7)} t}\right] + T_{38}^0 e^{-(b_{38}')^{(7)} t} \le T_{38}(t) \le$$

$$\frac{(a_{38})^{(7)} T_{36}^0}{(\mu_2)^{(7)}((R_1)^{(7)} + (r_{36})^{(7)} + (R_2)^{(7)})}\left[e^{((R_1)^{(7)} + (r_{36})^{(7)})t} - e^{-(R_2)^{(7)} t}\right] + T_{38}^0 e^{-(R_2)^{(7)} t}$$

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| **Definition of** $(S_1)^{(7)}, (S_2)^{(7)}, (R_1)^{(7)}, (R_2)^{(7)}$:-<br><br>Where $(S_1)^{(7)} = (a_{36})^{(7)}(m_2)^{(7)} - (a'_{36})^{(7)}$<br><br>$(S_2)^{(7)} = (a_{38})^{(7)} - (p_{38})^{(7)}$<br><br>$(R_1)^{(7)} = (b_{36})^{(7)}(\mu_2)^{(7)} - (b'_{36})^{(7)}$<br><br>$(R_2)^{(7)} = (b'_{38})^{(7)} - (r_{38})^{(7)}$ | 370 |
| **Behavior of the solutions of equation**<br><br>**Theorem 2:** If we denote and define<br><br>**Definition of** $(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ :<br><br>$(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ four constants satisfying<br><br>$-(\sigma_2)^{(8)} \le -(a'_{40})^{(8)} + (a'_{41})^{(8)} - (a''_{40})^{(8)}(T_{41}, t) + (a''_{41})^{(8)}(T_{41}, t) \le -(\sigma_1)^{(8)}$<br><br>$-(\tau_2)^{(8)} \le -(b'_{40})^{(8)} + (b'_{41})^{(8)} - (b''_{40})^{(8)}\big((G_{43}), t\big) - (b''_{41})^{(8)}\big((G_{43}), t\big) \le -(\tau_1)^{(8)}$ | 371 |
| **Definition of** $(v_1)^{(8)}, (v_2)^{(8)}, (u_1)^{(8)}, (u_2)^{(8)}, v^{(8)}, u^{(8)}$ :<br><br>By $(v_1)^{(8)} > 0, (v_2)^{(8)} < 0$ and respectively $(u_1)^{(8)} > 0, (u_2)^{(8)} < 0$ the roots of the equations $(a_{41})^{(8)}\big(v^{(8)}\big)^2 + (\sigma_1)^{(8)}v^{(8)} - (a_{40})^{(8)} = 0$<br>and $(b_{41})^{(8)}\big(u^{(8)}\big)^2 + (\tau_1)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$ and | 372 |
| **Definition of** $(\bar{v}_1)^{(8)}, , (\bar{v}_2)^{(8)}, (\bar{u}_1)^{(8)}, (\bar{u}_2)^{(8)}$ :<br><br>By $(\bar{v}_1)^{(8)} > 0, (\bar{v}_2)^{(8)} < 0$ and respectively $(\bar{u}_1)^{(8)} > 0, (\bar{u}_2)^{(8)} < 0$ the<br><br>roots of the equations $(a_{41})^{(8)}\big(v^{(8)}\big)^2 + (\sigma_2)^{(8)}v^{(8)} - (a_{40})^{(8)} = 0$<br><br>and $(b_{41})^{(8)}\big(u^{(8)}\big)^2 + (\tau_2)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$<br><br>**Definition of** $(m_1)^{(8)}, (m_2)^{(8)}, (\mu_1)^{(8)}, (\mu_2)^{(8)}, (v_0)^{(8)}$ :-<br><br>If we define $(m_1)^{(8)}, (m_2)^{(8)}, (\mu_1)^{(8)}, (\mu_2)^{(8)}$ by<br><br>$(m_2)^{(8)} = (v_0)^{(8)}, (m_1)^{(8)} = (v_1)^{(8)}$, ***if*** $(v_0)^{(8)} < (v_1)^{(8)}$<br><br>$(m_2)^{(8)} = (v_1)^{(8)}, (m_1)^{(8)} = (\bar{v}_1)^{(8)}$, ***if*** $(v_1)^{(8)} < (v_0)^{(8)} < (\bar{v}_1)^{(8)}$, | |

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| and $\boxed{(v_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}}$<br><br>$(m_2)^{(8)} = (v_1)^{(8)}, (m_1)^{(8)} = (v_0)^{(8)},$ ***if*** $(\bar{v}_1)^{(8)} < (v_0)^{(8)}$ | |
| and analogously<br><br>$(\mu_2)^{(8)} = (u_0)^{(8)}, (\mu_1)^{(8)} = (u_1)^{(8)},$ ***if*** $(u_0)^{(8)} < (u_1)^{(8)}$<br><br>$(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (\bar{u}_1)^{(8)}$ , ***if*** $(u_1)^{(8)} < (u_0)^{(8)} < (\bar{u}_1)^{(8)},$<br>and $\boxed{(u_0)^{(8)} = \frac{T_{40}^0}{T_{41}^0}}$<br><br>$(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (u_0)^{(8)},$ ***if*** $(\bar{u}_1)^{(8)} < (u_0)^{(8)}$ where $(u_1)^{(8)}, (\bar{u}_1)^{(8)}$ | 374 |
| Then the solution of global equations satisfies the inequalities<br><br>$G_{40}^0 e^{((S_1)^{(8)} - (p_{40})^{(8)})t} \leq G_{40}(t) \leq G_{40}^0 e^{(S_1)^{(8)} t}$<br><br>where $(p_i)^{(8)}$ is defined by equation | 375 |
| $\frac{1}{(m_1)^{(8)}} G_{40}^0 e^{((S_1)^{(8)} - (p_{40})^{(8)})t} \leq G_{41}(t) \leq \frac{1}{(m_2)^{(8)}} G_{40}^0 e^{(S_1)^{(8)} t}$ | 376 |
| $\left(\frac{(a_{42})^{(8)} G_{40}^0}{(m_1)^{(8)}((S_1)^{(8)} - (p_{40})^{(8)} - (S_2)^{(8)})}\left[e^{((S_1)^{(8)} - (p_{40})^{(8)})t} - e^{-(S_2)^{(8)} t}\right] + G_{42}^0 e^{-(S_2)^{(8)} t} \leq G_{42}(t) \leq \right.$<br>$\left.\frac{(a_{42})^{(8)} G_{40}^0}{(m_2)^{(8)}((S_1)^{(8)} - (a'_{42})^{(8)})}\left[e^{(S_1)^{(8)} t} - e^{-(a'_{42})^{(8)} t}\right] + G_{42}^0 e^{-(a'_{42})^{(8)} t}\right)$ | 377 |
| $\boxed{T_{40}^0 e^{(R_1)^{(8)} t} \leq T_{40}(t) \leq T_{40}^0 e^{((R_1)^{(8)} + (r_{40})^{(8)})t}}$ | 378 |
| $\frac{1}{(\mu_1)^{(8)}} T_{40}^0 e^{(R_1)^{(8)} t} \leq T_{40}(t) \leq \frac{1}{(\mu_2)^{(8)}} T_{40}^0 e^{((R_1)^{(8)} + (r_{40})^{(8)})t}$ | 379 |
| $\frac{(b_{42})^{(8)} T_{40}^0}{(\mu_1)^{(8)}((R_1)^{(8)} - (b'_{42})^{(8)})}\left[e^{(R_1)^{(8)} t} - e^{-(b'_{42})^{(8)} t}\right] + T_{42}^0 e^{-(b'_{42})^{(8)} t} \leq T_{42}(t) \leq$<br><br>$\frac{(a_{42})^{(8)} T_{40}^0}{(\mu_2)^{(8)}((R_1)^{(8)} + (r_{40})^{(8)} + (R_2)^{(8)})}\left[e^{((R_1)^{(8)} + (r_{40})^{(8)})t} - e^{-(R_2)^{(8)} t}\right] + T_{42}^0 e^{-(R_2)^{(8)} t}$ | 380 |
| **Definition of** $(S_1)^{(8)}, (S_2)^{(8)}, (R_1)^{(8)}, (R_2)^{(8)}$:-<br><br>Where $(S_1)^{(8)} = (a_{40})^{(8)}(m_2)^{(8)} - (a'_{40})^{(8)}$<br><br>$(S_2)^{(8)} = (a_{42})^{(8)} - (p_{42})^{(8)}$ | 381 |

$(R_1)^{(8)} = (b_{40})^{(8)}(\mu_2)^{(8)} - (b'_{40})^{(8)}$

$(R_2)^{(8)} = (b'_{42})^{(8)} - (r_{42})^{(8)}$

**Behavior of the solutions of equation 37 to 92** 382

**Theorem 2:** If we denote and define

**Definition of** $(\sigma_1)^{(9)}, (\sigma_2)^{(9)}, (\tau_1)^{(9)}, (\tau_2)^{(9)}$ :

$(\sigma_1)^{(9)}, (\sigma_2)^{(9)}, (\tau_1)^{(9)}, (\tau_2)^{(9)}$ four constants satisfying

$-(\sigma_2)^{(9)} \le -(a'_{44})^{(9)} + (a'_{45})^{(9)} - (a''_{44})^{(9)}(T_{45}, t) + (a''_{45})^{(9)}(T_{45}, t) \le -(\sigma_1)^{(9)}$

$-(\tau_2)^{(9)} \le -(b'_{44})^{(9)} + (b'_{45})^{(9)} - (b''_{44})^{(9)}\big((G_{47}), t\big) - (b''_{45})^{(9)}\big((G_{47}), t\big) \le -(\tau_1)^{(9)}$

**Definition of** $(\nu_1)^{(9)}, (\nu_2)^{(9)}, (u_1)^{(9)}, (u_2)^{(9)}, \nu^{(9)}, u^{(9)}$ :

By $(\nu_1)^{(9)} > 0$ , $(\nu_2)^{(9)} < 0$ and respectively $(u_1)^{(9)} > 0$ , $(u_2)^{(9)} < 0$ the roots of the equations $(a_{45})^{(9)}\big(\nu^{(9)}\big)^2 + (\sigma_1)^{(9)}\nu^{(9)} - (a_{44})^{(9)} = 0$ and $(b_{45})^{(9)}\big(u^{(9)}\big)^2 + (\tau_1)^{(9)}u^{(9)} - (b_{44})^{(9)} = 0$ and

**Definition of** $(\bar{\nu}_1)^{(9)}, , (\bar{\nu}_2)^{(9)}, (\bar{u}_1)^{(9)}, (\bar{u}_2)^{(9)}$ :

By $(\bar{\nu}_1)^{(9)} > 0$ , $(\bar{\nu}_2)^{(9)} < 0$ and respectively $(\bar{u}_1)^{(9)} > 0$ , $(\bar{u}_2)^{(9)} < 0$ the roots of the equations $(a_{45})^{(9)}\big(\nu^{(9)}\big)^2 + (\sigma_2)^{(9)}\nu^{(9)} - (a_{44})^{(9)} = 0$ and $(b_{45})^{(9)}\big(u^{(9)}\big)^2 + (\tau_2)^{(9)}u^{(9)} - (b_{44})^{(9)} = 0$

**Definition of** $(m_1)^{(9)}, (m_2)^{(9)}, (\mu_1)^{(9)}, (\mu_2)^{(9)}, (\nu_0)^{(9)}$ :-

If we define $(m_1)^{(9)}, (m_2)^{(9)}, (\mu_1)^{(9)}, (\mu_2)^{(9)}$ by

$(m_2)^{(9)} = (\nu_0)^{(9)}, (m_1)^{(9)} = (\nu_1)^{(9)}$, ***if*** $(\nu_0)^{(9)} < (\nu_1)^{(9)}$

$(m_2)^{(9)} = (\nu_1)^{(9)}, (m_1)^{(9)} = (\bar{\nu}_1)^{(9)}$ , ***if*** $(\nu_1)^{(9)} < (\nu_0)^{(9)} < (\bar{\nu}_1)^{(9)}$,

and $\boxed{(\nu_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}}}$

$(m_2)^{(9)} = (\nu_1)^{(9)}, (m_1)^{(9)} = (\nu_0)^{(9)}$, ***if*** $(\bar{\nu}_1)^{(9)} < (\nu_0)^{(9)}$

and analogously

$(\mu_2)^{(9)} = (u_0)^{(9)}, (\mu_1)^{(9)} = (u_1)^{(9)}$, ***if*** $(u_0)^{(9)} < (u_1)^{(9)}$

$(\mu_2)^{(9)} = (u_1)^{(9)}, (\mu_1)^{(9)} = (\bar{u}_1)^{(9)}$ , ***if*** $(u_1)^{(9)} < (u_0)^{(9)} < (\bar{u}_1)^{(9)}$,

and $\boxed{(u_0)^{(9)} = \frac{T^0_{44}}{T^0_{45}}}$

$(\mu_2)^{(9)} = (u_1)^{(9)}, (\mu_1)^{(9)} = (u_0)^{(9)}$, ***if*** $(\bar{u}_1)^{(9)} < (u_0)^{(9)}$ where $(u_1)^{(9)}, (\bar{u}_1)^{(9)}$ are defined by 59 and 69 respectively

Then the solution of 19,20,21,22,23 and 24 satisfies the inequalities

$$G_{44}^{0}e^{((S_1)^{(9)}-(p_{44})^{(9)})t} \le G_{44}(t) \le G_{44}^{0}e^{(S_1)^{(9)}t}$$

where $(p_i)^{(9)}$ is defined by equation 45

$$\frac{1}{(m_9)^{(9)}}G_{44}^{0}e^{((S_1)^{(9)}-(p_{44})^{(9)})t} \le G_{45}(t) \le \frac{1}{(m_2)^{(9)}}G_{44}^{0}e^{(S_1)^{(9)}t}$$

(

$$\frac{(a_{46})^{(9)}G_{44}^{0}}{(m_1)^{(9)}((S_1)^{(9)}-(p_{44})^{(9)}-(S_2)^{(9)})}\left[e^{((S_1)^{(9)}-(p_{44})^{(9)})t}-e^{-(S_2)^{(9)}t}\right]+G_{46}^{0}e^{-(S_2)^{(9)}t} \le G_{46}(t) \le$$

$$\frac{(a_{46})^{(9)}G_{44}^{0}}{(m_2)^{(9)}((S_1)^{(9)}-(a_{46}^{'})^{(9)})}[e^{(S_1)^{(9)}t}-e^{-(a_{46}^{'})^{(9)}t}]+\ G_{46}^{0}e^{-(a_{46}^{'})^{(9)}t})$$

$$\boxed{T_{44}^{0}e^{(R_1)^{(9)}t} \le T_{44}(t) \le T_{44}^{0}e^{((R_1)^{(9)}+(r_{44})^{(9)})t}}$$

$$\frac{1}{(\mu_1)^{(9)}}T_{44}^{0}e^{(R_1)^{(9)}t} \le T_{44}(t) \le \frac{1}{(\mu_2)^{(9)}}T_{44}^{0}e^{((R_1)^{(9)}+(r_{44})^{(9)})t}$$

$$\frac{(b_{46})^{(9)}T_{44}^{0}}{(\mu_1)^{(9)}((R_1)^{(9)}-(b_{46}^{'})^{(9)})}\left[e^{(R_1)^{(9)}t}-e^{-(b_{46}^{'})^{(9)}t}\right]+T_{46}^{0}e^{-(b_{46}^{'})^{(9)}t} \le T_{46}(t) \le$$

$$\frac{(a_{46})^{(9)}T_{44}^{0}}{(\mu_2)^{(9)}((R_1)^{(9)}+(r_{44})^{(9)}+(R_2)^{(9)})}\left[e^{((R_1)^{(9)}+(r_{44})^{(9)})t}-e^{-(R_2)^{(9)}t}\right]+T_{46}^{0}e^{-(R_2)^{(9)}t}$$

**Definition of** $(S_1)^{(9)}, (S_2)^{(9)}, (R_1)^{(9)}, (R_2)^{(9)}$:-

Where $(S_1)^{(9)} = (a_{44})^{(9)}(m_2)^{(9)} - (a_{44}^{'})^{(9)}$

$(S_2)^{(9)} = (a_{46})^{(9)} - (p_{46})^{(9)}$

$(R_1)^{(9)} = (b_{44})^{(9)}(\mu_2)^{(9)} - (b_{44}^{'})^{(9)}$

$(R_2)^{(9)} = (b_{46}^{'})^{(9)} - (r_{46})^{(9)}$

**Proof :** From global equations we obtain

$$\frac{dv^{(1)}}{dt} = (a_{13})^{(1)} - \left((a_{13}^{'})^{(1)} - (a_{14}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14},t)\right) - (a_{14}^{''})^{(1)}(T_{14},t)v^{(1)} - (a_{14})^{(1)}v^{(1)}$$

**Definition of** $v^{(1)}$ :- $\boxed{v^{(1)} = \frac{G_{13}}{G_{14}}}$

It follows

$$-\left((a_{14})^{(1)}(v^{(1)})^{2} + (\sigma_2)^{(1)}v^{(1)} - (a_{13})^{(1)}\right) \le \frac{dv^{(1)}}{dt} \le -\left((a_{14})^{(1)}(v^{(1)})^{2} + (\sigma_1)^{(1)}v^{(1)} - (a_{13})^{(1)}\right)$$

From which one obtains

**Definition of** $(\bar{v}_1)^{(1)}, (v_0)^{(1)}$ :-

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For $0 < \boxed{(\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}} < (\nu_1)^{(1)} < (\bar{\nu}_1)^{(1)}$

$$\nu^{(1)}(t) \geq \frac{(\nu_1)^{(1)} + (C)^{(1)}(\nu_2)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)} - (\nu_0)^{(1)}\right)t\right]}}{1 + (C)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)} - (\nu_0)^{(1)}\right)t\right]}} \quad , \quad \boxed{(C)^{(1)} = \frac{(\nu_1)^{(1)} - (\nu_0)^{(1)}}{(\nu_0)^{(1)} - (\nu_2)^{(1)}}}$$

it follows $(\nu_0)^{(1)} \leq \nu^{(1)}(t) \leq (\nu_1)^{(1)}$

In the same manner , we get 384

$$\nu^{(1)}(t) \leq \frac{(\bar{\nu}_1)^{(1)} + (\bar{C})^{(1)}(\bar{\nu}_2)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}}{1 + (\bar{C})^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(1)} = \frac{(\bar{\nu}_1)^{(1)} - (\nu_0)^{(1)}}{(\nu_0)^{(1)} - (\bar{\nu}_2)^{(1)}}}$$

From which we deduce $(\nu_0)^{(1)} \leq \nu^{(1)}(t) \leq (\bar{\nu}_1)^{(1)}$

If $0 < (\nu_1)^{(1)} < (\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0} < (\bar{\nu}_1)^{(1)}$ we find like in the previous case, 385

$$(\nu_1)^{(1)} \leq \frac{(\nu_1)^{(1)} + (C)^{(1)}(\nu_2)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)} - (\nu_2)^{(1)}\right)t\right]}}{1 + (C)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)} - (\nu_2)^{(1)}\right)t\right]}} \leq \nu^{(1)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(1)} + (\bar{C})^{(1)}(\bar{\nu}_2)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}}{1 + (\bar{C})^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}} \leq (\bar{\nu}_1)^{(1)}$$

If $0 < (\nu_1)^{(1)} \leq (\bar{\nu}_1)^{(1)} \leq \boxed{(\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}}$ , we obtain 386

$$(\nu_1)^{(1)} \leq \nu^{(1)}(t) \leq \frac{(\bar{\nu}_1)^{(1)} + (\bar{C})^{(1)}(\bar{\nu}_2)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}}{1 + (\bar{C})^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}} \leq (\nu_0)^{(1)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(1)}(t)$ :-

$$(m_2)^{(1)} \leq \nu^{(1)}(t) \leq (m_1)^{(1)}, \quad \boxed{\nu^{(1)}(t) = \frac{G_{13}(t)}{G_{14}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(1)}(t)$ :-

$$(\mu_2)^{(1)} \leq u^{(1)}(t) \leq (\mu_1)^{(1)}, \quad \boxed{u^{(1)}(t) = \frac{T_{13}(t)}{T_{14}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{13}^{''})^{(1)} = (a_{14}^{''})^{(1)}, then\ (\sigma_1)^{(1)} = (\sigma_2)^{(1)}$ and in this case $(\nu_1)^{(1)} = (\bar{\nu}_1)^{(1)}$ if in addition $(\nu_0)^{(1)} = (\nu_1)^{(1)}$ then $\nu^{(1)}(t) = (\nu_0)^{(1)}$ and as a consequence $G_{13}(t) = (\nu_0)^{(1)} G_{14}(t)$ this also defines $(\nu_0)^{(1)}$ for the special case

Analogously if $(b_{13}^{''})^{(1)} = (b_{14}^{''})^{(1)}, then\ (\tau_1)^{(1)} = (\tau_2)^{(1)}$ and then

$(u_1)^{(1)} = (\bar{u}_1)^{(1)}$ if in addition $(u_0)^{(1)} = (u_1)^{(1)}$ then $T_{13}(t) = (u_0)^{(1)} T_{14}(t)$ This is an important consequence of the relation between $(\nu_1)^{(1)}$ and $(\bar{\nu}_1)^{(1)}$, and definition of $(u_0)^{(1)}$.

**Proof** : From global equations we obtain 387

$$\frac{d\nu^{(2)}}{dt} = (a_{16})^{(2)} - \left((a_{16}^{'})^{(2)} - (a_{17}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17}, t)\right) - (a_{17}^{''})^{(2)}(T_{17}, t)\nu^{(2)} - (a_{17})^{(2)}\nu^{(2)}$$

**Definition of** $\nu^{(2)}$ :- $\boxed{\nu^{(2)} = \frac{G_{16}}{G_{17}}}$ 388

It follows 389

$$-\left((a_{17})^{(2)}\left(\nu^{(2)}\right)^2 + (\sigma_2)^{(2)}\nu^{(2)} - (a_{16})^{(2)}\right) \le \frac{d\nu^{(2)}}{dt} \le -\left((a_{17})^{(2)}\left(\nu^{(2)}\right)^2 + (\sigma_1)^{(2)}\nu^{(2)} - (a_{16})^{(2)}\right)$$

From which one obtains 390

**Definition of** $(\bar{\nu}_1)^{(2)}, (\nu_0)^{(2)}$ :-

For $0 < (\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0} < (\nu_1)^{(2)} < (\bar{\nu}_1)^{(2)}$

$$\nu^{(2)}(t) \ge \frac{(\nu_1)^{(2)} + (C)^{(2)}(\nu_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_0)^{(2)}\right)t\right]}}{1 + (C)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_0)^{(2)}\right)t\right]}} \quad , \quad \boxed{(C)^{(2)} = \frac{(\nu_1)^{(2)} - (\nu_0)^{(2)}}{(\nu_0)^{(2)} - (\nu_2)^{(2)}}}$$

it follows $(\nu_0)^{(2)} \le \nu^{(2)}(t) \le (\nu_1)^{(2)}$

In the same manner , we get 391

$$\nu^{(2)}(t) \le \frac{(\bar{\nu}_1)^{(2)} + (\bar{C})^{(2)}(\bar{\nu}_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}}{1 + (\bar{C})^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(2)} = \frac{(\bar{\nu}_1)^{(2)} - (\nu_0)^{(2)}}{(\nu_0)^{(2)} - (\bar{\nu}_2)^{(2)}}}$$

From which we deduce $(\nu_0)^{(2)} \le \nu^{(2)}(t) \le (\bar{\nu}_1)^{(2)}$ 392

If $0 < (\nu_1)^{(2)} < (\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0} < (\bar{\nu}_1)^{(2)}$ we find like in the previous case, 393

$$(\nu_1)^{(2)} \le \frac{(\nu_1)^{(2)} + (C)^{(2)}(\nu_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_2)^{(2)}\right)t\right]}}{1 + (C)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_2)^{(2)}\right)t\right]}} \le \nu^{(2)}(t) \le$$

| | |
|---|---|
| $\frac{(\bar{\nu}_1)^{(2)}+(\bar{C})^{(2)}(\bar{\nu}_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)}-(\bar{\nu}_2)^{(2)}\right)t\right]}}{1+(\bar{C})^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)}-(\bar{\nu}_2)^{(2)}\right)t\right]}} \leq (\bar{\nu}_1)^{(2)}$ | |
| If $0 < (\nu_1)^{(2)} \leq (\bar{\nu}_1)^{(2)} \leq (\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0}$ , we obtain<br><br>$(\nu_1)^{(2)} \leq \nu^{(2)}(t) \leq \frac{(\bar{\nu}_1)^{(2)}+(\bar{C})^{(2)}(\bar{\nu}_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)}-(\bar{\nu}_2)^{(2)}\right)t\right]}}{1+(\bar{C})^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)}-(\bar{\nu}_2)^{(2)}\right)t\right]}} \leq (\nu_0)^{(2)}$<br><br>And so with the notation of the first part of condition (c) , we have | 394 |
| **Definition of** $\nu^{(2)}(t)$ :-<br><br>$(m_2)^{(2)} \leq \nu^{(2)}(t) \leq (m_1)^{(2)}$, $\boxed{\nu^{(2)}(t) = \frac{G_{16}(t)}{G_{17}(t)}}$ | 395 |
| In a completely analogous way, we obtain<br><br>**Definition of** $u^{(2)}(t)$ :-<br><br>$(\mu_2)^{(2)} \leq u^{(2)}(t) \leq (\mu_1)^{(2)}$, $\boxed{u^{(2)}(t) = \frac{T_{16}(t)}{T_{17}(t)}}$ | 396 |
| Now, using this result and replacing it in global equations we get easily the result stated in the theorem. | |
| **Particular case :**<br><br>If $(a_{16}^{''})^{(2)} = (a_{17}^{''})^{(2)}, then\ (\sigma_1)^{(2)} = (\sigma_2)^{(2)}$ and in this case $(\nu_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ if in addition $(\nu_0)^{(2)} = (\nu_1)^{(2)}$ then $\nu^{(2)}(t) = (\nu_0)^{(2)}$ and as a consequence $G_{16}(t) = (\nu_0)^{(2)}G_{17}(t)$<br><br>Analogously if $(b_{16}^{''})^{(2)} = (b_{17}^{''})^{(2)}, then\ (\tau_1)^{(2)} = (\tau_2)^{(2)}$ and then<br><br>$(u_1)^{(2)} = (\bar{u}_1)^{(2)}$if in addition $(u_0)^{(2)} = (u_1)^{(2)}$ then $T_{16}(t) = (u_0)^{(2)}T_{17}(t)$ This is an important consequence of the relation between $(\nu_1)^{(2)}$ and $(\bar{\nu}_1)^{(2)}$ | 397 |
| **Proof** : From global equations we obtain<br><br>$\frac{d\nu^{(3)}}{dt} = (a_{20})^{(3)} - \left((a_{20}^{'})^{(3)} - (a_{21}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21},t)\right) - (a_{21}^{''})^{(3)}(T_{21},t)\nu^{(3)} - (a_{21})^{(3)}\nu^{(3)}$ | 398 |
| **Definition of** $\nu^{(3)}$ :- $\boxed{\nu^{(3)} = \frac{G_{20}}{G_{21}}}$<br><br>It follows<br><br>$-\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_2)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right) \leq \frac{d\nu^{(3)}}{dt} \leq -\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_1)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right)$ | 399 |
| | 400 |

From which one obtains

For $0 < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\nu_1)^{(3)} < (\bar{\nu}_1)^{(3)}$

$$\nu^{(3)}(t) \geq \frac{(\nu_1)^{(3)}+(C)^{(3)}(\nu_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_0)^{(3)}\right)t\right]}}{1+(C)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_0)^{(3)}\right)t\right]}} \quad , \quad \boxed{(C)^{(3)} = \frac{(\nu_1)^{(3)}-(\nu_0)^{(3)}}{(\nu_0)^{(3)}-(\nu_2)^{(3)}}}$$

it follows $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\nu_1)^{(3)}$

In the same manner , we get

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$$\nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(3)} = \frac{(\bar{\nu}_1)^{(3)}-(\nu_0)^{(3)}}{(\nu_0)^{(3)}-(\bar{\nu}_2)^{(3)}}}$$

**Definition of** $(\bar{\nu}_1)^{(3)}$ :-

From which we deduce $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\bar{\nu}_1)^{(3)}$

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If $0 < (\nu_1)^{(3)} < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\bar{\nu}_1)^{(3)}$ we find like in the previous case,

$$(\nu_1)^{(3)} \leq \frac{(\nu_1)^{(3)}+(C)^{(3)}(\nu_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}}{1+(C)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}} \leq \nu^{(3)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\bar{\nu}_1)^{(3)}$$

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If $0 < (\nu_1)^{(3)} \leq (\bar{\nu}_1)^{(3)} \leq (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0}$ , we obtain

$$(\nu_1)^{(3)} \leq \nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\nu_0)^{(3)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(3)}(t)$ :-

$$(m_2)^{(3)} \leq \nu^{(3)}(t) \leq (m_1)^{(3)}, \quad \boxed{\nu^{(3)}(t) = \frac{G_{20}(t)}{G_{21}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(3)}(t)$ :-

$$(\mu_2)^{(3)} \leq u^{(3)}(t) \leq (\mu_1)^{(3)}, \quad \boxed{u^{(3)}(t) = \frac{T_{20}(t)}{T_{21}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{20}^{''})^{(3)} = (a_{21}^{''})^{(3)}, then\ (\sigma_1)^{(3)} = (\sigma_2)^{(3)}$ and in this case $(\nu_1)^{(3)} = (\bar{\nu}_1)^{(3)}$ if in addition $(\nu_0)^{(3)} = (\nu_1)^{(3)}$ then $\nu^{(3)}(t) = (\nu_0)^{(3)}$ and as a consequence $G_{20}(t) = (\nu_0)^{(3)} G_{21}(t)$

Analogously if $(b_{20}^{''})^{(3)} = (b_{21}^{''})^{(3)}, then\ (\tau_1)^{(3)} = (\tau_2)^{(3)}$ and then

$(u_1)^{(3)} = (\bar{u}_1)^{(3)}$if in addition $(u_0)^{(3)} = (u_1)^{(3)}$ then $T_{20}(t) = (u_0)^{(3)} T_{21}(t)$ This is an important consequence of the relation between $(\nu_1)^{(3)}$ and $(\bar{\nu}_1)^{(3)}$

**Proof** : From global equations we obtain

$$\frac{d\nu^{(4)}}{dt} = (a_{24})^{(4)} - \left((a_{24}^{'})^{(4)} - (a_{25}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25}, t)\right) - (a_{25}^{''})^{(4)}(T_{25}, t)\nu^{(4)} - (a_{25})^{(4)}\nu^{(4)}$$

**Definition of** $\nu^{(4)}$ :- $\boxed{\nu^{(4)} = \frac{G_{24}}{G_{25}}}$

It follows

$$-\left((a_{25})^{(4)}\left(\nu^{(4)}\right)^2 + (\sigma_2)^{(4)}\nu^{(4)} - (a_{24})^{(4)}\right) \leq \frac{d\nu^{(4)}}{dt} \leq -\left((a_{25})^{(4)}\left(\nu^{(4)}\right)^2 + (\sigma_4)^{(4)}\nu^{(4)} - (a_{24})^{(4)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(4)}, (\nu_0)^{(4)}$ :-

For $0 < \boxed{(\nu_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0}} < (\nu_1)^{(4)} < (\bar{\nu}_1)^{(4)}$

$$\nu^{(4)}(t) \geq \frac{(\nu_1)^{(4)} + (C)^{(4)}(\nu_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_0)^{(4)}\right)t\right]}}{4 + (C)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_0)^{(4)}\right)t\right]}} \quad , \quad \boxed{(C)^{(4)} = \frac{(\nu_1)^{(4)} - (\nu_0)^{(4)}}{(\nu_0)^{(4)} - (\nu_2)^{(4)}}}$$

it follows $(\nu_0)^{(4)} \leq \nu^{(4)}(t) \leq (\nu_1)^{(4)}$

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In the same manner , we get

$$\nu^{(4)}(t) \leq \frac{(\bar{\nu}_1)^{(4)} + (\bar{C})^{(4)}(\bar{\nu}_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}}{4 + (\bar{C})^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(4)} = \frac{(\bar{\nu}_1)^{(4)} - (\nu_0)^{(4)}}{(\nu_0)^{(4)} - (\bar{\nu}_2)^{(4)}}}$$

From which we deduce $(\nu_0)^{(4)} \leq \nu^{(4)}(t) \leq (\bar{\nu}_1)^{(4)}$

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If $0 < (\nu_1)^{(4)} < (\nu_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0} < (\bar{\nu}_1)^{(4)}$ we find like in the previous case,

$$(\nu_1)^{(4)} \leq \frac{(\nu_1)^{(4)} + (C)^{(4)}(\nu_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_2)^{(4)}\right)t\right]}}{1 + (C)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_2)^{(4)}\right)t\right]}} \leq \nu^{(4)}(t) \leq$$

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$$\frac{(\bar{\nu}_1)^{(4)}+(\bar{C})^{(4)}(\bar{\nu}_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)}-(\bar{\nu}_2)^{(4)}\right)t\right]}}{1+(\bar{C})^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)}-(\bar{\nu}_2)^{(4)}\right)t\right]}} \leq (\bar{\nu}_1)^{(4)}$$

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If $0 < (\nu_1)^{(4)} \leq (\bar{\nu}_1)^{(4)} \leq \boxed{(\nu_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0}}$ , we obtain

$$(\nu_1)^{(4)} \leq \nu^{(4)}(t) \leq \frac{(\bar{\nu}_1)^{(4)}+(\bar{C})^{(4)}(\bar{\nu}_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)}-(\bar{\nu}_2)^{(4)}\right)t\right]}}{1+(\bar{C})^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)}-(\bar{\nu}_2)^{(4)}\right)t\right]}} \leq (\nu_0)^{(4)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(4)}(t)$ :-

$(m_2)^{(4)} \leq \nu^{(4)}(t) \leq (m_1)^{(4)}$, $\boxed{\nu^{(4)}(t) = \frac{G_{24}(t)}{G_{25}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(4)}(t)$ :-

$(\mu_2)^{(4)} \leq u^{(4)}(t) \leq (\mu_1)^{(4)}$, $\boxed{u^{(4)}(t) = \frac{T_{24}(t)}{T_{25}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{24}^{''})^{(4)} = (a_{25}^{''})^{(4)}, then\ (\sigma_1)^{(4)} = (\sigma_2)^{(4)}$ and in this case $(\nu_1)^{(4)} = (\bar{\nu}_1)^{(4)}$ if in addition $(\nu_0)^{(4)} = (\nu_1)^{(4)}$ then $\nu^{(4)}(t) = (\nu_0)^{(4)}$ and as a consequence $G_{24}(t) = (\nu_0)^{(4)}G_{25}(t)$ **this also defines** $(\nu_0)^{(4)}$ **for the special case** .

Analogously if $(b_{24}^{''})^{(4)} = (b_{25}^{''})^{(4)}, then\ (\tau_1)^{(4)} = (\tau_2)^{(4)}$ and then $(u_1)^{(4)} = (\bar{u}_4)^{(4)}$ if in addition $(u_0)^{(4)} = (u_1)^{(4)}$ then $T_{24}(t) = (u_0)^{(4)}T_{25}(t)$ This is an important consequence of the relation between $(\nu_1)^{(4)}$ and $(\bar{\nu}_1)^{(4)}$, **and definition of** $(u_0)^{(4)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(5)}}{dt} = (a_{28})^{(5)} - \left((a_{28}^{'})^{(5)} - (a_{29}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}, t)\right) - (a_{29}^{''})^{(5)}(T_{29}, t)\nu^{(5)} - (a_{29})^{(5)}\nu^{(5)}$$

**Definition of** $\nu^{(5)}$ :- $\boxed{\nu^{(5)} = \frac{G_{28}}{G_{29}}}$

It follows

$$-\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_2)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right) \leq \frac{d\nu^{(5)}}{dt} \leq -\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_1)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(5)}, (\nu_0)^{(5)}$ :-

For $0 < \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}} < (\nu_1)^{(5)} < (\bar{\nu}_1)^{(5)}$

$$\nu^{(5)}(t) \geq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}}{5 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}} \quad , \quad \boxed{(C)^{(5)} = \frac{(\nu_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\nu_2)^{(5)}}}$$

it follows $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\nu_1)^{(5)}$

In the same manner , we get

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$$\nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)} + (\bar{C})^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{5 + (\bar{C})^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(5)} = \frac{(\bar{\nu}_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\bar{\nu}_2)^{(5)}}}$$

From which we deduce $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\bar{\nu}_5)^{(5)}$

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If $0 < (\nu_1)^{(5)} < (\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0} < (\bar{\nu}_1)^{(5)}$ we find like in the previous case,

$$(\nu_1)^{(5)} \leq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}}{1 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}} \leq \nu^{(5)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(5)} + (C)^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{1 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\bar{\nu}_1)^{(5)}$$

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If $0 < (\nu_1)^{(5)} \leq (\bar{\nu}_1)^{(5)} \leq \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}}$ , we obtain

$$(\nu_1)^{(5)} \leq \nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)} + (C)^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{1 + (\bar{C})^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\nu_0)^{(5)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(5)}(t)$ :-

$$(m_2)^{(5)} \leq \nu^{(5)}(t) \leq (m_1)^{(5)}, \quad \boxed{\nu^{(5)}(t) = \frac{G_{28}(t)}{G_{29}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(5)}(t)$ :-

$$(\mu_2)^{(5)} \leq u^{(5)}(t) \leq (\mu_1)^{(5)}, \quad \boxed{u^{(5)}(t) = \frac{T_{28}(t)}{T_{29}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{28}'')^{(5)} = (a_{29}'')^{(5)}, then\ (\sigma_1)^{(5)} = (\sigma_2)^{(5)}$ and in this case $(\nu_1)^{(5)} = (\bar{\nu}_1)^{(5)}$ if in addition $(\nu_0)^{(5)} = (\nu_5)^{(5)}$ then $\nu^{(5)}(t) = (\nu_0)^{(5)}$ and as a consequence $G_{28}(t) = (\nu_0)^{(5)} G_{29}(t)$ **this also defines** $(\nu_0)^{(5)}$ **for the special case** .

Analogously if $(b_{28}'')^{(5)} = (b_{29}'')^{(5)}, then\ (\tau_1)^{(5)} = (\tau_2)^{(5)}$ and then

$(u_1)^{(5)} = (\bar{u}_1)^{(5)}$if in addition $(u_0)^{(5)} = (u_1)^{(5)}$ then $T_{28}(t) = (u_0)^{(5)}T_{29}(t)$ This is an important consequence of the relation between $(\nu_1)^{(5)}$ and $(\bar{\nu}_1)^{(5)}$, **and definition of** $(u_0)^{(5)}$.

**Proof** : From global equations we obtain

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$$\frac{d\nu^{(6)}}{dt} = (a_{32})^{(6)} - \left((a'_{32})^{(6)} - (a'_{33})^{(6)} + (a''_{32})^{(6)}(T_{33},t)\right) - (a''_{33})^{(6)}(T_{33},t)\nu^{(6)} - (a_{33})^{(6)}\nu^{(6)}$$

**Definition of** $\nu^{(6)}$ :- $\boxed{\nu^{(6)} = \frac{G_{32}}{G_{33}}}$

It follows

$$-\left((a_{33})^{(6)}\left(\nu^{(6)}\right)^2 + (\sigma_2)^{(6)}\nu^{(6)} - (a_{32})^{(6)}\right) \le \frac{d\nu^{(6)}}{dt} \le -\left((a_{33})^{(6)}\left(\nu^{(6)}\right)^2 + (\sigma_1)^{(6)}\nu^{(6)} - (a_{32})^{(6)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(6)}, (\nu_0)^{(6)}$ :-

For $0 < \boxed{(\nu_0)^{(6)} = \frac{G^0_{32}}{G^0_{33}}} < (\nu_1)^{(6)} < (\bar{\nu}_1)^{(6)}$

$$\nu^{(6)}(t) \ge \frac{(\nu_1)^{(6)} + (C)^{(6)}(\nu_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)} - (\nu_0)^{(6)}\right)t\right]}}{1 + (C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)} - (\nu_0)^{(6)}\right)t\right]}} \quad , \quad \boxed{(C)^{(6)} = \frac{(\nu_1)^{(6)} - (\nu_0)^{(6)}}{(\nu_0)^{(6)} - (\nu_2)^{(6)}}}$$

it follows $(\nu_0)^{(6)} \le \nu^{(6)}(t) \le (\nu_1)^{(6)}$

In the same manner , we get

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$$\nu^{(6)}(t) \le \frac{(\bar{\nu}_1)^{(6)} + (\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)} - (\bar{\nu}_2)^{(6)}\right)t\right]}}{1 + (\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)} - (\bar{\nu}_2)^{(6)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(6)} = \frac{(\bar{\nu}_1)^{(6)} - (\nu_0)^{(6)}}{(\nu_0)^{(6)} - (\bar{\nu}_2)^{(6)}}}$$

From which we deduce $(\nu_0)^{(6)} \le \nu^{(6)}(t) \le (\bar{\nu}_1)^{(6)}$

If $0 < (\nu_1)^{(6)} < (\nu_0)^{(6)} = \frac{G^0_{32}}{G^0_{33}} < (\bar{\nu}_1)^{(6)}$ we find like in the previous case,

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$$(\nu_1)^{(6)} \le \frac{(\nu_1)^{(6)} + (C)^{(6)}(\nu_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)} - (\nu_2)^{(6)}\right)t\right]}}{1 + (C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)} - (\nu_2)^{(6)}\right)t\right]}} \le \nu^{(6)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(6)} + (\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)} - (\bar{\nu}_2)^{(6)}\right)t\right]}}{1 + (\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)} - (\bar{\nu}_2)^{(6)}\right)t\right]}} \le (\bar{\nu}_1)^{(6)}$$

If $0 < (\nu_1)^{(6)} \le (\bar{\nu}_1)^{(6)} \le \boxed{(\nu_0)^{(6)} = \frac{G^0_{32}}{G^0_{33}}}$ , we obtain

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$$(\nu_1)^{(6)} \le \nu^{(6)}(t) \le \frac{(\bar{\nu}_1)^{(6)} + (\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)} - (\bar{\nu}_2)^{(6)}\right)t\right]}}{1 + (C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)} - (\bar{\nu}_2)^{(6)}\right)t\right]}} \le (\nu_0)^{(6)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(6)}(t)$ :-

$(m_2)^{(6)} \leq \nu^{(6)}(t) \leq (m_1)^{(6)}$, $\boxed{\nu^{(6)}(t) = \frac{G_{32}(t)}{G_{33}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(6)}(t)$ :-

$(\mu_2)^{(6)} \leq u^{(6)}(t) \leq (\mu_1)^{(6)}$, $\boxed{u^{(6)}(t) = \frac{T_{32}(t)}{T_{33}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{32}^{''})^{(6)} = (a_{33}^{''})^{(6)}, then\ (\sigma_1)^{(6)} = (\sigma_2)^{(6)}$ and in this case $(\nu_1)^{(6)} = (\bar{\nu}_1)^{(6)}$ if in addition $(\nu_0)^{(6)} = (\nu_1)^{(6)}$ then $\nu^{(6)}(t) = (\nu_0)^{(6)}$ and as a consequence $G_{32}(t) = (\nu_0)^{(6)} G_{33}(t)$ **this also defines** $(\nu_0)^{(6)}$ **for the special case** .

Analogously if $(b_{32}^{''})^{(6)} = (b_{33}^{''})^{(6)}, then\ (\tau_1)^{(6)} = (\tau_2)^{(6)}$ and then $(u_1)^{(6)} = (\bar{u}_1)^{(6)}$ if in addition $(u_0)^{(6)} = (u_1)^{(6)}$ then $T_{32}(t) = (u_0)^{(6)} T_{33}(t)$ This is an important consequence of the relation between $(\nu_1)^{(6)}$ and $(\bar{\nu}_1)^{(6)}$, **and definition of** $(u_0)^{(6)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(7)}}{dt} = (a_{36})^{(7)} - \left((a_{36}^{'})^{(7)} - (a_{37}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37}, t)\right) - (a_{37}^{''})^{(7)}(T_{37}, t)\nu^{(7)} - (a_{37})^{(7)}\nu^{(7)}$$

**Definition of** $\nu^{(7)}$ :- $\boxed{\nu^{(7)} = \frac{G_{36}}{G_{37}}}$

It follows

$$-\left((a_{37})^{(7)}\left(\nu^{(7)}\right)^2 + (\sigma_2)^{(7)}\nu^{(7)} - (a_{36})^{(7)}\right) \leq \frac{d\nu^{(7)}}{dt} \leq -\left((a_{37})^{(7)}\left(\nu^{(7)}\right)^2 + (\sigma_1)^{(7)}\nu^{(7)} - (a_{36})^{(7)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(7)}, (\nu_0)^{(7)}$ :-

For $0 < \boxed{(\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0}} < (\nu_1)^{(7)} < (\bar{\nu}_1)^{(7)}$

$$\nu^{(7)}(t) \geq \frac{(\nu_1)^{(7)} + (C)^{(7)}(\nu_2)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)} - (\nu_0)^{(7)}\right)t\right]}}{1 + (C)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)} - (\nu_0)^{(7)}\right)t\right]}} \quad , \quad \boxed{(C)^{(7)} = \frac{(\nu_1)^{(7)} - (\nu_0)^{(7)}}{(\nu_0)^{(7)} - (\nu_2)^{(7)}}}$$

it follows $(\nu_0)^{(7)} \leq \nu^{(7)}(t) \leq (\nu_1)^{(7)}$

In the same manner , we get

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$$\nu^{(7)}(t) \le \frac{(\bar{\nu}_1)^{(7)}+(\bar{C})^{(7)}(\bar{\nu}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}}{1+(\bar{C})^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(7)} = \frac{(\bar{\nu}_1)^{(7)}-(\nu_0)^{(7)}}{(\nu_0)^{(7)}-(\bar{\nu}_2)^{(7)}}}$$

From which we deduce $(\nu_0)^{(7)} \le \nu^{(7)}(t) \le (\bar{\nu}_1)^{(7)}$

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If $0 < (\nu_1)^{(7)} < (\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0} < (\bar{\nu}_1)^{(7)}$ we find like in the previous case,

$$(\nu_1)^{(7)} \le \frac{(\nu_1)^{(7)}+(C)^{(7)}(\nu_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)}-(\nu_2)^{(7)}\right)t\right]}}{1+(C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)}-(\nu_2)^{(7)}\right)t\right]}} \le \nu^{(7)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(7)}+(\bar{C})^{(7)}(\bar{\nu}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}}{1+(\bar{C})^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}} \le (\bar{\nu}_1)^{(7)}$$

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If $0 < (\nu_1)^{(7)} \le (\bar{\nu}_1)^{(7)} \le \boxed{(\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0}}$ , we obtain

$$(\nu_1)^{(7)} \le \nu^{(7)}(t) \le \frac{(\bar{\nu}_1)^{(7)}+(C)^{(7)}(\bar{\nu}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}}{1+(C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}} \le (\nu_0)^{(7)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(7)}(t)$ :-

$$(m_2)^{(7)} \le \nu^{(7)}(t) \le (m_1)^{(7)}, \quad \boxed{\nu^{(7)}(t) = \frac{G_{36}(t)}{G_{37}(t)}}$$

In a completely analogous way, we obtain

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**Definition of** $u^{(7)}(t)$ :-

$$(\mu_2)^{(7)} \le u^{(7)}(t) \le (\mu_1)^{(7)}, \quad \boxed{u^{(7)}(t) = \frac{T_{36}(t)}{T_{37}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{36}'')^{(7)} = (a_{37}'')^{(7)}, then\ (\sigma_1)^{(7)} = (\sigma_2)^{(7)}$ and in this case $(\nu_1)^{(7)} = (\bar{\nu}_1)^{(7)}$ if in addition $(\nu_0)^{(7)} = (\nu_1)^{(7)}$ then $\nu^{(7)}(t) = (\nu_0)^{(7)}$ and as a consequence $G_{36}(t) = (\nu_0)^{(7)}G_{37}(t)$ **this also defines** $(\nu_0)^{(7)}$ **for the special case** .

Analogously if $(b_{36}'')^{(7)} = (b_{37}'')^{(7)}, then\ (\tau_1)^{(7)} = (\tau_2)^{(7)}$ and then $(u_1)^{(7)} = (\bar{u}_1)^{(7)}$if in addition $(u_0)^{(7)} = (u_1)^{(7)}$ then $T_{36}(t) = (u_0)^{(7)}T_{37}(t)$ This is an important consequence of the relation between $(\nu_1)^{(7)}$ and $(\bar{\nu}_1)^{(7)}$, **and definition of** $(u_0)^{(7)}$.

**Proof** : From global equations we obtain

$$\frac{dv^{(8)}}{dt} = (a_{40})^{(8)} - \left((a_{40}^{'})^{(8)} - (a_{41}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41},t)\right) - (a_{41}^{''})^{(8)}(T_{41},t)v^{(8)} - (a_{41})^{(8)}v^{(8)}$$

**Definition of** $v^{(8)}$ :- $\boxed{v^{(8)} = \frac{G_{40}}{G_{41}}}$

It follows

$$-\left((a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_2)^{(8)}v^{(8)} - (a_{40})^{(8)}\right) \le \frac{dv^{(8)}}{dt} \le -\left((a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_1)^{(8)}v^{(8)} - (a_{40})^{(8)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(8)}, (\nu_0)^{(8)}$ :-

For $0 < \boxed{(\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}} < (\nu_1)^{(8)} < (\bar{\nu}_1)^{(8)}$

$$\nu^{(8)}(t) \ge \frac{(\nu_1)^{(8)} + (C)^{(8)}(\nu_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}}{1 + (C)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}} \quad , \quad \boxed{(C)^{(8)} = \frac{(\nu_1)^{(8)} - (\nu_0)^{(8)}}{(\nu_0)^{(8)} - (\nu_2)^{(8)}}}$$

it follows $(\nu_0)^{(8)} \le \nu^{(8)}(t) \le (\nu_1)^{(8)}$

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In the same manner , we get

$$\nu^{(8)}(t) \le \frac{(\bar{\nu}_1)^{(8)} + (\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}}{1 + (\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(8)} = \frac{(\bar{\nu}_1)^{(8)} - (\nu_0)^{(8)}}{(\nu_0)^{(8)} - (\bar{\nu}_2)^{(8)}}}$$

From which we deduce $(\nu_0)^{(8)} \le \nu^{(8)}(t) \le (\bar{\nu}_8)^{(8)}$

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If $0 < (\nu_1)^{(8)} < (\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0} < (\bar{\nu}_1)^{(8)}$ we find like in the previous case,

$$(\nu_1)^{(8)} \le \frac{(\nu_1)^{(8)} + (C)^{(8)}(\nu_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_2)^{(8)}\right)t\right]}}{1 + (C)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_2)^{(8)}\right)t\right]}} \le \nu^{(8)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(8)} + (C)^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}}{1 + (\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}} \le (\bar{\nu}_1)^{(8)}$$

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If $0 < (\nu_1)^{(8)} \le (\bar{\nu}_1)^{(8)} \le \boxed{(\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}}$ , we obtain

$$(\nu_1)^{(8)} \le \nu^{(8)}(t) \le \frac{(\bar{\nu}_1)^{(8)} + (\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}}{1 + (\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}} \le (\nu_0)^{(8)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $v^{(8)}(t)$ :-

$(m_2)^{(8)} \leq v^{(8)}(t) \leq (m_1)^{(8)}$, $\boxed{v^{(8)}(t) = \frac{G_{40}(t)}{G_{41}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(8)}(t)$ :-

$(\mu_2)^{(8)} \leq u^{(8)}(t) \leq (\mu_1)^{(8)}$, $\boxed{u^{(8)}(t) = \frac{T_{40}(t)}{T_{41}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{40}^{''})^{(8)} = (a_{41}^{''})^{(8)}, then\ (\sigma_1)^{(8)} = (\sigma_2)^{(8)}$ and in this case $(\nu_1)^{(8)} = (\bar{\nu}_1)^{(8)}$ if in addition $(\nu_0)^{(8)} = (\nu_1)^{(8)}$ then $v^{(8)}(t) = (\nu_0)^{(8)}$ and as a consequence $G_{40}(t) = (\nu_0)^{(8)} G_{41}(t)$ **this also defines** $(\nu_0)^{(8)}$ **for the special case** .

Analogously if $(b_{40}^{''})^{(8)} = (b_{41}^{''})^{(8)}, then\ (\tau_1)^{(8)} = (\tau_2)^{(8)}$ and then
$(u_1)^{(8)} = (\bar{u}_1)^{(8)}$if in addition $(u_0)^{(8)} = (u_1)^{(8)}$ then $T_{40}(t) = (u_0)^{(8)} T_{41}(t)$ This is an important consequence of the relation between $(\nu_1)^{(8)}$ and $(\bar{\nu}_1)^{(8)}$, **and definition of** $(u_0)^{(8)}$.

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**Proof** : From 99,20,44,22,23,44 we obtain

$$\frac{dv^{(9)}}{dt} = (a_{44})^{(9)} - \left((a_{44}^{'})^{(9)} - (a_{45}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45}, t)\right) - (a_{45}^{''})^{(9)}(T_{45}, t)\nu^{(9)} - (a_{45})^{(9)}\nu^{(9)}$$

**Definition of** $\nu^{(9)}$ :- $\boxed{\nu^{(9)} = \frac{G_{44}}{G_{45}}}$

It follows

$$-\left((a_{45})^{(9)}\left(\nu^{(9)}\right)^2 + (\sigma_2)^{(9)}\nu^{(9)} - (a_{44})^{(9)}\right) \leq \frac{d\nu^{(9)}}{dt} \leq -\left((a_{45})^{(9)}\left(\nu^{(9)}\right)^2 + (\sigma_1)^{(9)}\nu^{(9)} - (a_{44})^{(9)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(9)}, (\nu_0)^{(9)}$ :-

For $0 < \boxed{(\nu_0)^{(9)} = \frac{G_{44}^0}{G_{45}^0}} < (\nu_1)^{(9)} < (\bar{\nu}_1)^{(9)}$

$$v^{(9)}(t) \geq \frac{(\nu_1)^{(9)} + (C)^{(9)}(\nu_2)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_0)^{(9)}\right)t\right]}}{1 + (C)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_0)^{(9)}\right)t\right]}} \quad , \quad \boxed{(C)^{(9)} = \frac{(\nu_1)^{(9)} - (\nu_0)^{(9)}}{(\nu_0)^{(9)} - (\nu_2)^{(9)}}}$$

it follows $(\nu_0)^{(9)} \leq v^{(9)}(t) \leq (\nu_9)^{(9)}$

In the same manner , we get

$$\nu^{(9)}(t) \le \frac{(\bar{\nu}_1)^{(9)}+(\bar{C})^{(9)}(\bar{\nu}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)}-(\bar{\nu}_2)^{(9)}\right)t\right]}}{1+(\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)}-(\bar{\nu}_2)^{(9)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(9)} = \frac{(\bar{\nu}_1)^{(9)}-(\nu_0)^{(9)}}{(\nu_0)^{(9)}-(\bar{\nu}_2)^{(9)}}}$$

From which we deduce $(\nu_0)^{(9)} \le \nu^{(9)}(t) \le (\bar{\nu}_1)^{(9)}$

If $0 < (\nu_1)^{(9)} < (\nu_0)^{(9)} = \frac{G_{44}^0}{G_{45}^0} < (\bar{\nu}_1)^{(9)}$ we find like in the previous case,

$$(\nu_1)^{(9)} \le \frac{(\nu_1)^{(9)}+(C)^{(9)}(\nu_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)}-(\nu_2)^{(9)}\right)t\right]}}{1+(C)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)}-(\nu_2)^{(9)}\right)t\right]}} \le \nu^{(9)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(9)}+(\bar{C})^{(9)}(\bar{\nu}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)}-(\bar{\nu}_2)^{(9)}\right)t\right]}}{1+(\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)}-(\bar{\nu}_2)^{(9)}\right)t\right]}} \le (\bar{\nu}_1)^{(9)}$$

If $0 < (\nu_1)^{(9)} \le (\bar{\nu}_1)^{(9)} \le \boxed{(\nu_0)^{(9)} = \frac{G_{44}^0}{G_{45}^0}}$, we obtain

$$(\nu_1)^{(9)} \le \nu^{(9)}(t) \le \frac{(\bar{\nu}_1)^{(9)}+(\bar{C})^{(9)}(\bar{\nu}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)}-(\bar{\nu}_2)^{(9)}\right)t\right]}}{1+(\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)}-(\bar{\nu}_2)^{(9)}\right)t\right]}} \le (\nu_0)^{(9)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(9)}(t)$ :-

$$(m_2)^{(9)} \le \nu^{(9)}(t) \le (m_1)^{(9)}, \quad \boxed{\nu^{(9)}(t) = \frac{G_{44}(t)}{G_{45}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(9)}(t)$ :-

$$(\mu_2)^{(9)} \le u^{(9)}(t) \le (\mu_1)^{(9)}, \quad \boxed{u^{(9)}(t) = \frac{T_{44}(t)}{T_{45}(t)}}$$

Now, using this result and replacing it in 99, 20,44,22,23, and 44 we get easily the result stated in the theorem.

**Particular case :**

If $(a_{44}^{''})^{(9)} = (a_{45}^{''})^{(9)}, then\ (\sigma_1)^{(9)} = (\sigma_2)^{(9)}$ and in this case $(\nu_1)^{(9)} = (\bar{\nu}_1)^{(9)}$ if in addition $(\nu_0)^{(9)} = (\nu_1)^{(9)}$ then $\nu^{(9)}(t) = (\nu_0)^{(9)}$ and as a consequence $G_{44}(t) = (\nu_0)^{(9)}G_{45}(t)$ **this also defines** $(\nu_0)^{(9)}$ **for the special case** .

Analogously if $(b_{44}^{''})^{(9)} = (b_{45}^{''})^{(9)}, then\ (\tau_1)^{(9)} = (\tau_2)^{(9)}$ and then $(u_1)^{(9)} = (\bar{u}_1)^{(9)}$ if in addition $(u_0)^{(9)} = (u_1)^{(9)}$ then $T_{44}(t) = (u_0)^{(9)}T_{45}(t)$ This is an important consequence of the relation between $(\nu_1)^{(9)}$ and $(\bar{\nu}_1)^{(9)}$, **and definition of** $(u_0)^{(9)}$.

We can prove the following

**Theorem :** If $(a_i^{''})^{(1)} and\ (b_i^{''})^{(1)}$ are independent on $t$ , and the conditions with the notations

$$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} < 0$$

$$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a_{13})^{(1)}(p_{13})^{(1)} + (a_{14}^{'})^{(1)}(p_{14})^{(1)} + (p_{13})^{(1)}(p_{14})^{(1)} > 0$$

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| $(b_{13}')^{(1)}(b_{14}')^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} > 0$ ,<br>$(b_{13}')^{(1)}(b_{14}')^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} - (b_{13}')^{(1)}(r_{14})^{(1)} - (b_{14}')^{(1)}(r_{14})^{(1)} + (r_{13})^{(1)}(r_{14})^{(1)} < 0$<br>$with\ \ (p_{13})^{(1)}, (r_{14})^{(1)}$ as defined by equation are satisfied , then the system | |
| **Theorem :** If $(a_i'')^{(2)} and\ (b_i'')^{(2)}$ are independent on t , and the conditions with the notations | 426 |
| $(a_{16}')^{(2)}(a_{17}')^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} < 0$ | 427 |
| $(a_{16}')^{(2)}(a_{17}')^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a_{16})^{(2)}(p_{16})^{(2)} + (a_{17}')^{(2)}(p_{17})^{(2)} + (p_{16})^{(2)}(p_{17})^{(2)} > 0$ | 428 |
| $(b_{16}')^{(2)}(b_{17}')^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} > 0$ , | 429 |
| $(b_{16}')^{(2)}(b_{17}')^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} - (b_{16}')^{(2)}(r_{17})^{(2)} - (b_{17}')^{(2)}(r_{17})^{(2)} + (r_{16})^{(2)}(r_{17})^{(2)} < 0$<br>$with\ \ (p_{16})^{(2)}, (r_{17})^{(2)}$ as defined by equation are satisfied , then the system | 430 |
| **Theorem :** If $(a_i'')^{(3)} and\ (b_i'')^{(3)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{20}')^{(3)}(a_{21}')^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} < 0$<br>$(a_{20}')^{(3)}(a_{21}')^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a_{20})^{(3)}(p_{20})^{(3)} + (a_{21}')^{(3)}(p_{21})^{(3)} + (p_{20})^{(3)}(p_{21})^{(3)} > 0$<br>$(b_{20}')^{(3)}(b_{21}')^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} > 0$ ,<br>$(b_{20}')^{(3)}(b_{21}')^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} - (b_{20}')^{(3)}(r_{21})^{(3)} - (b_{21}')^{(3)}(r_{21})^{(3)} + (r_{20})^{(3)}(r_{21})^{(3)} < 0$<br>$with\ \ (p_{20})^{(3)}, (r_{21})^{(3)}$ as defined by equation are satisfied , then the system | 431 |
| We can prove the following<br>**Theorem :** If $(a_i'')^{(4)} and\ (b_i'')^{(4)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{24}')^{(4)}(a_{25}')^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} < 0$<br>$(a_{24}')^{(4)}(a_{25}')^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a_{24})^{(4)}(p_{24})^{(4)} + (a_{25}')^{(4)}(p_{25})^{(4)} + (p_{24})^{(4)}(p_{25})^{(4)} > 0$<br>$(b_{24}')^{(4)}(b_{25}')^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} > 0$ ,<br>$(b_{24}')^{(4)}(b_{25}')^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} - (b_{24}')^{(4)}(r_{25})^{(4)} - (b_{25}')^{(4)}(r_{25})^{(4)} + (r_{24})^{(4)}(r_{25})^{(4)} < 0$<br>$with\ \ (p_{24})^{(4)}, (r_{25})^{(4)}$ as defined by equation are satisfied , then the system | 432 |
| **Theorem :** If $(a_i'')^{(5)} and\ (b_i'')^{(5)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{28}')^{(5)}(a_{29}')^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} < 0$<br>$(a_{28}')^{(5)}(a_{29}')^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a_{28})^{(5)}(p_{28})^{(5)} + (a_{29}')^{(5)}(p_{29})^{(5)} + (p_{28})^{(5)}(p_{29})^{(5)} > 0$ | 433 |

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| $(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} > 0$ ,<br>$(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} - (b_{28}^{'})^{(5)}(r_{29})^{(5)} - (b_{29}^{'})^{(5)}(r_{29})^{(5)} + (r_{28})^{(5)}(r_{29})^{(5)} < 0$<br>$with\ \ (p_{28})^{(5)}, (r_{29})^{(5)}$ as defined by equation are satisfied , then the system | |
| **Theorem** If $(a_i^{''})^{(6)} and\ (b_i^{''})^{(6)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} < 0$<br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32})^{(6)}(p_{32})^{(6)} + (a_{33}^{'})^{(6)}(p_{33})^{(6)} + (p_{32})^{(6)}(p_{33})^{(6)} > 0$<br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} > 0$ ,<br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} - (b_{32}^{'})^{(6)}(r_{33})^{(6)} - (b_{33}^{'})^{(6)}(r_{33})^{(6)} + (r_{32})^{(6)}(r_{33})^{(6)} < 0$<br>$with\ \ (p_{32})^{(6)}, (r_{33})^{(6)}$ as defined by equation are satisfied , then the system | 434 |
| **Theorem :** If $(a_i^{''})^{(7)} and\ (b_i^{''})^{(7)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} < 0$<br>$(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36})^{(7)}(p_{36})^{(7)} + (a_{37}^{'})^{(7)}(p_{37})^{(7)} + (p_{36})^{(7)}(p_{37})^{(7)} > 0$<br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} > 0$ ,<br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} - (b_{36}^{'})^{(7)}(r_{37})^{(7)} - (b_{37}^{'})^{(7)}(r_{37})^{(7)} + (r_{36})^{(7)}(r_{37})^{(7)} < 0$<br>$with\ \ (p_{36})^{(7)}, (r_{37})^{(7)}$ as defined by equation are satisfied , then the system | 435 |
| **Theorem :** If $(a_i^{''})^{(8)} and\ (b_i^{''})^{(8)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} < 0$<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40})^{(8)}(p_{40})^{(8)} + (a_{41}^{'})^{(8)}(p_{41})^{(8)} + (p_{40})^{(8)}(p_{41})^{(8)} > 0$<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} > 0$ ,<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} - (b_{40}^{'})^{(8)}(r_{41})^{(8)} - (b_{41}^{'})^{(8)}(r_{41})^{(8)} + (r_{40})^{(8)}(r_{41})^{(8)} < 0$<br>$with\ \ (p_{40})^{(8)}, (r_{41})^{(8)}$ as defined by equation are satisfied , then the system | 436 |
| **Theorem :** If $(a_i^{''})^{(9)} and\ (b_i^{''})^{(9)}$ are independent on $t$ , and the conditions (with the notations 45,46,27,28)<br>$(a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} < 0$ | 436 A |

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| $(a'_{44})^{(9)}(a'_{45})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44})^{(9)}(p_{44})^{(9)} + (a'_{45})^{(9)}(p_{45})^{(9)} + (p_{44})^{(9)}(p_{45})^{(9)} > 0$<br>$(b'_{44})^{(9)}(b'_{45})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} > 0$ ,<br>$(b'_{44})^{(9)}(b'_{45})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} - (b'_{44})^{(9)}(r_{45})^{(9)} - (b'_{45})^{(9)}(r_{45})^{(9)} + (r_{44})^{(9)}(r_{45})^{(9)} < 0$<br>$with\ \ (p_{44})^{(9)}, (r_{45})^{(9)}$ as defined by equation 45 are satisfied , then the system | |
| $(a_{13})^{(1)}G_{14} - \left[(a'_{13})^{(1)} + (a''_{13})^{(1)}(T_{14})\right]G_{13} = 0$ | 437 |
| $(a_{14})^{(1)}G_{13} - \left[(a'_{14})^{(1)} + (a''_{14})^{(1)}(T_{14})\right]G_{14} = 0$ | 438 |
| $(a_{15})^{(1)}G_{14} - \left[(a'_{15})^{(1)} + (a''_{15})^{(1)}(T_{14})\right]G_{15} = 0$ | 439 |
| $(b_{13})^{(1)}T_{14} - [(b'_{13})^{(1)} - (b''_{13})^{(1)}(G)\ ]T_{13} = 0$ | 440 |
| $(b_{14})^{(1)}T_{13} - [(b'_{14})^{(1)} - (b''_{14})^{(1)}(G)\ ]T_{14} = 0$ | 441 |
| $(b_{15})^{(1)}T_{14} - [(b'_{15})^{(1)} - (b''_{15})^{(1)}(G)\ ]T_{15} = 0$ | 442 |
| has a unique positive solution , which is an equilibrium solution for the system | |
| $(a_{16})^{(2)}G_{17} - \left[(a'_{16})^{(2)} + (a''_{16})^{(2)}(T_{17})\right]G_{16} = 0$ | 443 |
| $(a_{17})^{(2)}G_{16} - \left[(a'_{17})^{(2)} + (a''_{17})^{(2)}(T_{17})\right]G_{17} = 0$ | 444 |
| $(a_{18})^{(2)}G_{17} - \left[(a'_{18})^{(2)} + (a''_{18})^{(2)}(T_{17})\right]G_{18} = 0$ | 445 |
| $(b_{16})^{(2)}T_{17} - [(b'_{16})^{(2)} - (b''_{16})^{(2)}(G_{19})\ ]T_{16} = 0$ | 446 |
| $(b_{17})^{(2)}T_{16} - [(b'_{17})^{(2)} - (b''_{17})^{(2)}(G_{19})\ ]T_{17} = 0$ | 447 |
| $(b_{18})^{(2)}T_{17} - [(b'_{18})^{(2)} - (b''_{18})^{(2)}(G_{19})\ ]T_{18} = 0$ | 448 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{20})^{(3)}G_{21} - \left[(a'_{20})^{(3)} + (a''_{20})^{(3)}(T_{21})\right]G_{20} = 0$ | 449 |
| $(a_{21})^{(3)}G_{20} - \left[(a'_{21})^{(3)} + (a''_{21})^{(3)}(T_{21})\right]G_{21} = 0$ | 450 |
| $(a_{22})^{(3)}G_{21} - \left[(a'_{22})^{(3)} + (a''_{22})^{(3)}(T_{21})\right]G_{22} = 0$ | 451 |
| $(b_{20})^{(3)}T_{21} - [(b'_{20})^{(3)} - (b''_{20})^{(3)}(G_{23})\ ]T_{20} = 0$ | 452 |
| $(b_{21})^{(3)}T_{20} - [(b'_{21})^{(3)} - (b''_{21})^{(3)}(G_{23})\ ]T_{21} = 0$ | 453 |
| $(b_{22})^{(3)}T_{21} - [(b'_{22})^{(3)} - (b''_{22})^{(3)}(G_{23})\ ]T_{22} = 0$ | 454 |
| has a unique positive solution , which is an equilibrium solution | |

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| $(a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25})\right]G_{24} = 0$ | 455 |
| $(a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25})\right]G_{25} = 0$ | 456 |
| $(a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25})\right]G_{26} = 0$ | 457 |
| $(b_{24})^{(4)}T_{25} - [(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}\left((G_{27})\right)]T_{24} = 0$ | 458 |
| $(b_{25})^{(4)}T_{24} - [(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}\left((G_{27})\right)]T_{25} = 0$ | 459 |
| $(b_{26})^{(4)}T_{25} - [(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}\left((G_{27})\right)]T_{26} = 0$ | 460 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29})\right]G_{28} = 0$ | 461 |
| $(a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29})\right]G_{29} = 0$ | 462 |
| $(a_{30})^{(5)}G_{29} - \left[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29})\right]G_{30} = 0$ | 463 |
| $(b_{28})^{(5)}T_{29} - [(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}(G_{31})]T_{28} = 0$ | 464 |
| $(b_{29})^{(5)}T_{28} - [(b_{29}^{'})^{(5)} - (b_{29}^{''})^{(5)}(G_{31})]T_{29} = 0$ | 465 |
| $(b_{30})^{(5)}T_{29} - [(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}(G_{31})]T_{30} = 0$ | 466 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{32})^{(6)}G_{33} - \left[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33})\right]G_{32} = 0$ | 467 |
| $(a_{33})^{(6)}G_{32} - \left[(a_{33}^{'})^{(6)} + (a_{33}^{''})^{(6)}(T_{33})\right]G_{33} = 0$ | 468 |
| $(a_{34})^{(6)}G_{33} - \left[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33})\right]G_{34} = 0$ | 469 |
| $(b_{32})^{(6)}T_{33} - [(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}(G_{35})]T_{32} = 0$ | 470 |
| $(b_{33})^{(6)}T_{32} - [(b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}(G_{35})]T_{33} = 0$ | 471 |
| $(b_{34})^{(6)}T_{33} - [(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}(G_{35})]T_{34} = 0$ | 472 |
| has a unique positive solution , which is an equilibrium solution | |

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| $(a_{36})^{(7)}G_{37} - \left[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37})\right]G_{36} = 0$ | 473 |
| $(a_{37})^{(7)}G_{36} - \left[(a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37})\right]G_{37} = 0$ | 474 |
| $(a_{38})^{(7)}G_{37} - \left[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37})\right]G_{38} = 0$ | 475 |
| $(b_{36})^{(7)}T_{37} - [(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}(G_{39})]T_{36} = 0$ | 476 |
| $(b_{37})^{(7)}T_{36} - [(b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}(G_{39})\,]T_{37} = 0$ | 477 |
| $(b_{38})^{(7)}T_{37} - [(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}(G_{39})\,]T_{38} = 0$ | 478 |
| $(a_{40})^{(8)}G_{41} - \left[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41})\right]G_{40} = 0$ | 479 |
| $(a_{41})^{(8)}G_{40} - \left[(a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41})\right]G_{41} = 0$ | 480 |
| $(a_{42})^{(8)}G_{41} - \left[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41})\right]G_{42} = 0$ | 481 |
| $(b_{40})^{(8)}T_{41} - [(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}(G_{43})\,]T_{40} = 0$ | 482 |
| $(b_{41})^{(8)}T_{40} - [(b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}(G_{43})\,]T_{41} = 0$ | 483 |
| $(b_{42})^{(8)}T_{41} - [(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}(G_{43})\,]T_{42} = 0$ | 484 |
| $(a_{44})^{(9)}G_{45} - \left[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45})\right]G_{44} = 0$ | 484 A |
| $(a_{45})^{(9)}G_{44} - \left[(a_{45}^{'})^{(9)} + (a_{45}^{''})^{(9)}(T_{45})\right]G_{45} = 0$ | |
| $(a_{46})^{(9)}G_{45} - \left[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45})\right]G_{46} = 0$ | |
| $(b_{44})^{(9)}T_{45} - [(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}(G_{47})\,]T_{44} = 0$ | |
| $(b_{45})^{(9)}T_{44} - [(b_{45}^{'})^{(9)} - (b_{45}^{''})^{(9)}(G_{47})\,]T_{45} = 0$ | |
| $(b_{46})^{(9)}T_{45} - [(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}(G_{47})\,]T_{46} = 0$ | |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{13}, G_{14}$ if<br><br>$F(T) = (a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a_{13}^{'})^{(1)}(a_{14}^{''})^{(1)}(T_{14}) + (a_{14}^{'})^{(1)}(a_{13}^{''})^{(1)}(T_{14}) +$ | 485 |

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| $(a_{13}^{''})^{(1)}(T_{14})(a_{14}^{''})^{(1)}(T_{14}) = 0$ | |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{16}, G_{17}$ if<br><br>$\mathrm{F}(T_{19}) = (a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a_{16}^{'})^{(2)}(a_{17}^{''})^{(2)}(T_{17}) + (a_{17}^{'})^{(2)}(a_{16}^{''})^{(2)}(T_{17}) +$<br>$(a_{16}^{''})^{(2)}(T_{17})(a_{17}^{''})^{(2)}(T_{17}) = 0$ | 486 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{20}, G_{21}$ if<br><br>$F(T_{23}) = (a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a_{20}^{'})^{(3)}(a_{21}^{''})^{(3)}(T_{21}) + (a_{21}^{'})^{(3)}(a_{20}^{''})^{(3)}(T_{21}) +$<br>$(a_{20}^{''})^{(3)}(T_{21})(a_{21}^{''})^{(3)}(T_{21}) = 0$ | 487 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{24}, G_{25}$ if<br><br>$F(T_{27}) = (a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a_{24}^{'})^{(4)}(a_{25}^{''})^{(4)}(T_{25}) + (a_{25}^{'})^{(4)}(a_{24}^{''})^{(4)}(T_{25}) +$<br>$(a_{24}^{''})^{(4)}(T_{25})(a_{25}^{''})^{(4)}(T_{25}) = 0$ | 488 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{28}, G_{29}$ if<br><br>$F(T_{31}) = (a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a_{28}^{'})^{(5)}(a_{29}^{''})^{(5)}(T_{29}) + (a_{29}^{'})^{(5)}(a_{28}^{''})^{(5)}(T_{29}) +$<br>$(a_{28}^{''})^{(5)}(T_{29})(a_{29}^{''})^{(5)}(T_{29}) = 0$ | 489 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{32}, G_{33}$ if<br><br>$F(T_{35}) = (a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32}^{'})^{(6)}(a_{33}^{''})^{(6)}(T_{33}) + (a_{33}^{'})^{(6)}(a_{32}^{''})^{(6)}(T_{33}) +$<br>$(a_{32}^{''})^{(6)}(T_{33})(a_{33}^{''})^{(6)}(T_{33}) = 0$ | 490 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{36}, G_{37}$ if<br><br>$F(T_{39}) = (a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36}^{'})^{(7)}(a_{37}^{''})^{(7)}(T_{37}) + (a_{37}^{'})^{(7)}(a_{36}^{''})^{(7)}(T_{37}) +$<br>$(a_{36}^{''})^{(7)}(T_{37})(a_{37}^{''})^{(7)}(T_{37}) = 0$ | 491 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{40}, G_{41}$ if<br><br>$F(T_{43}) = (a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40}^{'})^{(8)}(a_{41}^{''})^{(8)}(T_{41}) + (a_{41}^{'})^{(8)}(a_{40}^{''})^{(8)}(T_{41}) +$<br>$(a_{40}^{''})^{(8)}(T_{41})(a_{41}^{''})^{(8)}(T_{41}) = 0$ | 492 |

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| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{44}, G_{45}$ if<br><br>$F(T_{47}) = (a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44}^{'})^{(9)}(a_{45}^{''})^{(9)}(T_{45}) + (a_{45}^{'})^{(9)}(a_{44}^{''})^{(9)}(T_{45}) + (a_{44}^{''})^{(9)}(T_{45})(a_{45}^{''})^{(9)}(T_{45}) = 0$ | 492<br>A |
| **Definition and uniqueness of** $T_{14}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{14})$ being increasing, it follows that there exists a unique $T_{14}^{*}$ for which $f(T_{14}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{13} = \frac{(a_{13})^{(1)}G_{14}}{[(a_{13}^{'})^{(1)}+(a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15} = \frac{(a_{15})^{(1)}G_{14}}{[(a_{15}^{'})^{(1)}+(a_{15}^{''})^{(1)}(T_{14}^{*})]}$ | 493 |
| **Definition and uniqueness of** $T_{17}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(2)}(T_{17})$ being increasing, it follows that there exists a unique $T_{17}^{*}$ for which $f(T_{17}^{*}) = 0$. With this value , we obtain from the three first equations | 494 |
| $G_{16} = \frac{(a_{16})^{(2)}G_{17}}{[(a_{16}^{'})^{(2)}+(a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18} = \frac{(a_{18})^{(2)}G_{17}}{[(a_{18}^{'})^{(2)}+(a_{18}^{''})^{(2)}(T_{17}^{*})]}$ | 495 |
| **Definition and uniqueness of** $T_{21}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{21})$ being increasing, it follows that there exists a unique $T_{21}^{*}$ for which $f(T_{21}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{20} = \frac{(a_{20})^{(3)}G_{21}}{[(a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21}^{*})]}$ , $G_{22} = \frac{(a_{22})^{(3)}G_{21}}{[(a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21}^{*})]}$ | 496 |
| **Definition and uniqueness of** $T_{25}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(4)}(T_{25})$ being increasing, it follows that there exists a unique $T_{25}^{*}$ for which $f(T_{25}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{24} = \frac{(a_{24})^{(4)}G_{25}}{[(a_{24}^{'})^{(4)}+(a_{24}^{''})^{(4)}(T_{25}^{*})]}$ , $G_{26} = \frac{(a_{26})^{(4)}G_{25}}{[(a_{26}^{'})^{(4)}+(a_{26}^{''})^{(4)}(T_{25}^{*})]}$ | 497 |
| **Definition and uniqueness of** $T_{29}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(5)}(T_{29})$ being increasing, it follows that there exists a unique $T_{29}^{*}$ for which $f(T_{29}^{*}) = 0$. With this value , we obtain from the three first equations | 498 |

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| $G_{28} = \frac{(a_{28})^{(5)}G_{29}}{[(a'_{28})^{(5)}+(a''_{28})^{(5)}(T^*_{29})]}$ , $G_{30} = \frac{(a_{30})^{(5)}G_{29}}{[(a'_{30})^{(5)}+(a''_{30})^{(5)}(T^*_{29})]}$ | |
| **Definition and uniqueness of** $T^*_{33}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a''_i)^{(6)}(T_{33})$ being increasing, it follows that there exists a unique $T^*_{33}$ for which $f(T^*_{33}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{32} = \frac{(a_{32})^{(6)}G_{33}}{[(a'_{32})^{(6)}+(a''_{32})^{(6)}(T^*_{33})]}$ , $G_{34} = \frac{(a_{34})^{(6)}G_{33}}{[(a'_{34})^{(6)}+(a''_{34})^{(6)}(T^*_{33})]}$ | 499 |
| **Definition and uniqueness of** $T^*_{37}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a''_i)^{(7)}(T_{37})$ being increasing, it follows that there exists a unique $T^*_{37}$ for which $f(T^*_{37}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{36} = \frac{(a_{36})^{(7)}G_{37}}{[(a'_{36})^{(7)}+(a''_{36})^{(7)}(T^*_{37})]}$ , $G_{38} = \frac{(a_{38})^{(7)}G_{37}}{[(a'_{38})^{(7)}+(a''_{38})^{(7)}(T^*_{37})]}$ | 500 |
| **Definition and uniqueness of** $T^*_{41}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a''_i)^{(8)}(T_{41})$ being increasing, it follows that there exists a unique $T^*_{41}$ for which $f(T^*_{41}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{40} = \frac{(a_{40})^{(8)}G_{41}}{[(a'_{40})^{(8)}+(a''_{40})^{(8)}(T^*_{41})]}$ , $G_{42} = \frac{(a_{42})^{(8)}G_{41}}{[(a'_{42})^{(8)}+(a''_{42})^{(8)}(T^*_{41})]}$ | 501 |
| **Definition and uniqueness of** $T^*_{45}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a''_i)^{(9)}(T_{45})$ being increasing, it follows that there exists a unique $T^*_{45}$ for which $f(T^*_{45}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{44} = \frac{(a_{44})^{(9)}G_{45}}{[(a'_{44})^{(9)}+(a''_{44})^{(9)}(T^*_{45})]}$ , $G_{46} = \frac{(a_{46})^{(9)}G_{45}}{[(a'_{46})^{(9)}+(a''_{46})^{(9)}(T^*_{45})]}$ | 501 A |
| By the same argument, the equations admit solutions $G_{13}, G_{14}$ if<br>$\varphi(G) = (b'_{13})^{(1)}(b'_{14})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} -$<br><br>$\left[(b'_{13})^{(1)}(b''_{14})^{(1)}(G) + (b'_{14})^{(1)}(b''_{13})^{(1)}(G)\right]+(b''_{13})^{(1)}(G)(b''_{14})^{(1)}(G) = 0$<br><br>Where in $G(G_{13}, G_{14}, G_{15}), G_{13}, G_{15}$ must be replaced by their values from 96. It is easy to see that $\varphi$ is a decreasing function in $G_{14}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that | 502 |

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| there exists a unique $G_{14}^{*}$ such that $\varphi(G^{*})=0$ | |
| By the same argument, the equations admit solutions $G_{16}, G_{17}$ if<br>$\varphi(G_{19})=(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)}-(b_{16})^{(2)}(b_{17})^{(2)}-$<br>$\left[(b_{16}^{'})^{(2)}(b_{17}^{''})^{(2)}(G_{19})+(b_{17}^{'})^{(2)}(b_{16}^{''})^{(2)}(G_{19})\right]+(b_{16}^{''})^{(2)}(G_{19})(b_{17}^{''})^{(2)}(G_{19})=0$ | 503 |
| Where in $(G_{19})(G_{16},G_{17},G_{18}), G_{16}, G_{18}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{17}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{14}^{*}$ such that $\varphi((G_{19})^{*})=0$ | 504 |
| By the same argument, the equations admit solutions $G_{20}, G_{21}$ if<br>$\varphi(G_{23})=(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)}-(b_{20})^{(3)}(b_{21})^{(3)}-$<br>$\left[(b_{20}^{'})^{(3)}(b_{21}^{''})^{(3)}(G_{23})+(b_{21}^{'})^{(3)}(b_{20}^{''})^{(3)}(G_{23})\right]+(b_{20}^{''})^{(3)}(G_{23})(b_{21}^{''})^{(3)}(G_{23})=0$<br>Where in $G_{23}(G_{20},G_{21},G_{22}), G_{20}, G_{22}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{21}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{21}^{*}$ such that $\varphi((G_{23})^{*})=0$ | 505 |
| By the same argument, the equations admit solutions $G_{24}, G_{25}$ if<br>$\varphi(G_{27})=(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)}-(b_{24})^{(4)}(b_{25})^{(4)}-$<br>$\left[(b_{24}^{'})^{(4)}(b_{25}^{''})^{(4)}(G_{27})+(b_{25}^{'})^{(4)}(b_{24}^{''})^{(4)}(G_{27})\right]+(b_{24}^{''})^{(4)}(G_{27})(b_{25}^{''})^{(4)}(G_{27})=0$<br>Where in $(G_{27})(G_{24},G_{25},G_{26}), G_{24}, G_{26}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{25}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{25}^{*}$ such that $\varphi((G_{27})^{*})=0$ | 506 |
| By the same argument, the equations admit solutions $G_{28}, G_{29}$ if<br>$\varphi(G_{31})=(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)}-(b_{28})^{(5)}(b_{29})^{(5)}-$<br>$\left[(b_{28}^{'})^{(5)}(b_{29}^{''})^{(5)}(G_{31})+(b_{29}^{'})^{(5)}(b_{28}^{''})^{(5)}(G_{31})\right]+(b_{28}^{''})^{(5)}(G_{31})(b_{29}^{''})^{(5)}(G_{31})=0$<br>Where in $(G_{31})(G_{28},G_{29},G_{30}), G_{28}, G_{30}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{29}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{29}^{*}$ such that $\varphi((G_{31})^{*})=0$ | 507 |
| By the same argument, the equations admit solutions $G_{32}, G_{33}$ if<br>$\varphi(G_{35})=(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)}-(b_{32})^{(6)}(b_{33})^{(6)}-$<br>$\left[(b_{32}^{'})^{(6)}(b_{33}^{''})^{(6)}(G_{35})+(b_{33}^{'})^{(6)}(b_{32}^{''})^{(6)}(G_{35})\right]+(b_{32}^{''})^{(6)}(G_{35})(b_{33}^{''})^{(6)}(G_{35})=0$<br>Where in $(G_{35})(G_{32},G_{33},G_{34}), G_{32}, G_{34}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{33}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that | 508 |

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| there exists a unique $G_{33}^{*}$ such that $\varphi(G_{35}{}^{*}) = 0$ | |
| By the same argument, the equations admit solutions $G_{36}, G_{37}$ if<br>$\varphi(G_{39}) = (b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} -$<br>$\left[(b_{36}^{'})^{(7)}(b_{37}^{''})^{(7)}(G_{39}) + (b_{37}^{'})^{(7)}(b_{36}^{''})^{(7)}(G_{39})\right] + (b_{36}^{''})^{(7)}(G_{39})(b_{37}^{''})^{(7)}(G_{39}) = 0$<br><br>Where in $(G_{39})(G_{36}, G_{37}, G_{38}), G_{36}, G_{38}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{37}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{37}^{*}$ such that $\varphi(G_{39}{}^{*}) = 0$ | 509 |
| By the same argument, the equations admit solutions $G_{40}, G_{41}$ if<br>$\varphi(G_{43}) = (b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} -$<br>$\left[(b_{40}^{'})^{(8)}(b_{41}^{''})^{(8)}(G_{43}) + (b_{41}^{'})^{(8)}(b_{40}^{''})^{(8)}(G_{43})\right] + (b_{40}^{''})^{(8)}(G_{43})(b_{41}^{''})^{(8)}(G_{43}) = 0$<br><br>Where in $(G_{43})(G_{40}, G_{41}, G_{42}), G_{40}, G_{42}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{41}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{41}^{*}$ such that $\varphi(G_{43}{}^{*}) = 0$ | 510 |
| By the same argument, the equations 92,93 admit solutions $G_{44}, G_{45}$ if<br>$\varphi(G_{47}) = (b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} -$<br><br>$\left[(b_{44}^{'})^{(9)}(b_{45}^{''})^{(9)}(G_{47}) + (b_{45}^{'})^{(9)}(b_{44}^{''})^{(9)}(G_{47})\right] + (b_{44}^{''})^{(9)}(G_{47})(b_{45}^{''})^{(9)}(G_{47}) = 0$<br><br>Where in $(G_{47})(G_{44}, G_{45}, G_{46}), G_{44}, G_{46}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{45}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{45}^{*}$ such that $\varphi((G_{47})^{*}) = 0$ | |
| Finally we obtain the unique solution<br><br>$G_{14}^{*}$ given by $\varphi(G^{*}) = 0$ , $T_{14}^{*}$ given by $f(T_{14}^{*}) = 0$ and<br><br>$G_{13}^{*} = \frac{(a_{13})^{(1)}G_{14}^{*}}{[(a_{13}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15}^{*} = \frac{(a_{15})^{(1)}G_{14}^{*}}{[(a_{15}^{'})^{(1)} + (a_{15}^{''})^{(1)}(T_{14}^{*})]}$<br><br>$T_{13}^{*} = \frac{(b_{13})^{(1)}T_{14}^{*}}{[(b_{13}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G^{*})]}$ , $T_{15}^{*} = \frac{(b_{15})^{(1)}T_{14}^{*}}{[(b_{15}^{'})^{(1)} - (b_{15}^{''})^{(1)}(G^{*})]}$<br><br>Obviously, these values represent an equilibrium solution | 511 |
| Finally we obtain the unique solution | |
| $G_{17}^{*}$ given by $\varphi((G_{19})^{*}) = 0$ , $T_{17}^{*}$ given by $f(T_{17}^{*}) = 0$ and | 512 |
| $G_{16}^{*} = \frac{(a_{16})^{(2)}G_{17}^{*}}{[(a_{16}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18}^{*} = \frac{(a_{18})^{(2)}G_{17}^{*}}{[(a_{18}^{'})^{(2)} + (a_{18}^{''})^{(2)}(T_{17}^{*})]}$ | 513 |
| $T_{16}^{*} = \frac{(b_{16})^{(2)}T_{17}^{*}}{[(b_{16}^{'})^{(2)} - (b_{16}^{''})^{(2)}((G_{19})^{*})]}$ , $T_{18}^{*} = \frac{(b_{18})^{(2)}T_{17}^{*}}{[(b_{18}^{'})^{(2)} - (b_{18}^{''})^{(2)}((G_{19})^{*})]}$ | 514 |

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| Obviously, these values represent an equilibrium solution | |
| Finally we obtain the unique solution<br><br>$G_{21}^{*}$ given by $\varphi((G_{23})^{*}) = 0$ , $T_{21}^{*}$ given by $f(T_{21}^{*}) = 0$ and<br><br>$G_{20}^{*} = \frac{(a_{20})^{(3)} G_{21}^{*}}{[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21}^{*})]}$ , $G_{22}^{*} = \frac{(a_{22})^{(3)} G_{21}^{*}}{[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21}^{*})]}$<br><br>$T_{20}^{*} = \frac{(b_{20})^{(3)} T_{21}^{*}}{[(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23}{}^{*})]}$ , $T_{22}^{*} = \frac{(b_{22})^{(3)} T_{21}^{*}}{[(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23}{}^{*})]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 515 |
| Finally we obtain the unique solution<br><br>$G_{25}^{*}$ given by $\varphi(G_{27}) = 0$ , $T_{25}^{*}$ given by $f(T_{25}^{*}) = 0$ and<br><br>$G_{24}^{*} = \frac{(a_{24})^{(4)} G_{25}^{*}}{[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25}^{*})]}$ , $G_{26}^{*} = \frac{(a_{26})^{(4)} G_{25}^{*}}{[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25}^{*})]}$ | 516 |
| $T_{24}^{*} = \frac{(b_{24})^{(4)} T_{25}^{*}}{[(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}((G_{27})^{*})]}$ , $T_{26}^{*} = \frac{(b_{26})^{(4)} T_{25}^{*}}{[(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}((G_{27})^{*})]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 517 |
| Finally we obtain the unique solution<br><br>$G_{29}^{*}$ given by $\varphi((G_{31})^{*}) = 0$ , $T_{29}^{*}$ given by $f(T_{29}^{*}) = 0$ and<br><br>$G_{28}^{*} = \frac{(a_{28})^{(5)} G_{29}^{*}}{[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}^{*})]}$ , $G_{30}^{*} = \frac{(a_{30})^{(5)} G_{29}^{*}}{[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29}^{*})]}$ | 518 |
| $T_{28}^{*} = \frac{(b_{28})^{(5)} T_{29}^{*}}{[(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}((G_{31})^{*})]}$ , $T_{30}^{*} = \frac{(b_{30})^{(5)} T_{29}^{*}}{[(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}((G_{31})^{*})]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 519 |
| Finally we obtain the unique solution<br><br>$G_{33}^{*}$ given by $\varphi((G_{35})^{*}) = 0$ , $T_{33}^{*}$ given by $f(T_{33}^{*}) = 0$ and<br><br>$G_{32}^{*} = \frac{(a_{32})^{(6)} G_{33}^{*}}{[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33}^{*})]}$ , $G_{34}^{*} = \frac{(a_{34})^{(6)} G_{33}^{*}}{[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33}^{*})]}$ | 520 |
| $T_{32}^{*} = \frac{(b_{32})^{(6)} T_{33}^{*}}{[(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}((G_{35})^{*})]}$ , $T_{34}^{*} = \frac{(b_{34})^{(6)} T_{33}^{*}}{[(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}((G_{35})^{*})]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 521 |
| Finally we obtain the unique solution<br><br>$G_{37}^{*}$ given by $\varphi((G_{39})^{*}) = 0$ , $T_{37}^{*}$ given by $f(T_{37}^{*}) = 0$ and | 522 |

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| $G_{36}^{*} = \frac{(a_{36})^{(7)} G_{37}^{*}}{[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)} (T_{37}^{*})]}$ , $G_{38}^{*} = \frac{(a_{38})^{(7)} G_{37}^{*}}{[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)} (T_{37}^{*})]}$<br><br>$T_{36}^{*} = \frac{(b_{36})^{(7)} T_{37}^{*}}{[(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)} ((G_{39})^{*})]}$ , $T_{38}^{*} = \frac{(b_{38})^{(7)} T_{37}^{*}}{[(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)} ((G_{39})^{*})]}$ | |
| Finally we obtain the unique solution<br><br>$G_{41}^{*}$ given by $\varphi((G_{43})^{*}) = 0$ , $T_{41}^{*}$ given by $f(T_{41}^{*}) = 0$ and<br><br>$G_{40}^{*} = \frac{(a_{40})^{(8)} G_{41}^{*}}{[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)} (T_{41}^{*})]}$ , $G_{42}^{*} = \frac{(a_{42})^{(8)} G_{41}^{*}}{[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)} (T_{41}^{*})]}$<br><br>$T_{40}^{*} = \frac{(b_{40})^{(8)} T_{41}^{*}}{[(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)} ((G_{43})^{*})]}$ , $T_{42}^{*} = \frac{(b_{42})^{(8)} T_{41}^{*}}{[(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)} ((G_{43})^{*})]}$ | 523 |
| Finally we obtain the unique solution of 89 to 99<br><br>$G_{45}^{*}$ given by $\varphi((G_{47})^{*}) = 0$ , $T_{45}^{*}$ given by $f(T_{45}^{*}) = 0$ and<br><br>$G_{44}^{*} = \frac{(a_{44})^{(9)} G_{45}^{*}}{[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)} (T_{45}^{*})]}$ , $G_{46}^{*} = \frac{(a_{46})^{(9)} G_{45}^{*}}{[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)} (T_{45}^{*})]}$<br><br>$T_{44}^{*} = \frac{(b_{44})^{(9)} T_{45}^{*}}{[(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)} ((G_{47})^{*})]}$ , $T_{46}^{*} = \frac{(b_{46})^{(9)} T_{45}^{*}}{[(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)} ((G_{47})^{*})]}$ | 523 A |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(1)}$ *and* $(b_i^{''})^{(1)}$ Belong to $C^{(1)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>**Proof:**_Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^{*} + \mathbb{G}_i$ , $T_i = T_i^{*} + \mathbb{T}_i$<br><br>$\frac{\partial (a_{14}^{''})^{(1)}}{\partial T_{14}} (T_{14}^{*}) = (q_{14})^{(1)}$ , $\frac{\partial (b_i^{''})^{(1)}}{\partial G_j} (G^{*}) = s_{ij}$ | 524 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{13}}{dt} = -\left((a_{13}^{'})^{(1)} + (p_{13})^{(1)}\right)\mathbb{G}_{13} + (a_{13})^{(1)}\mathbb{G}_{14} - (q_{13})^{(1)} G_{13}^{*} \mathbb{T}_{14}$ | 525 |

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| $\frac{d\mathbb{G}_{14}}{dt} = -\left((a_{14}^{'})^{(1)} + (p_{14})^{(1)}\right)\mathbb{G}_{14} + (a_{14})^{(1)}\mathbb{G}_{13} - (q_{14})^{(1)}G_{14}^{*}\mathbb{T}_{14}$ | 526 |
| $\frac{d\mathbb{G}_{15}}{dt} = -\left((a_{15}^{'})^{(1)} + (p_{15})^{(1)}\right)\mathbb{G}_{15} + (a_{15})^{(1)}\mathbb{G}_{14} - (q_{15})^{(1)}G_{15}^{*}\mathbb{T}_{14}$ | 527 |
| $\frac{d\mathbb{T}_{13}}{dt} = -\left((b_{13}^{'})^{(1)} - (r_{13})^{(1)}\right)\mathbb{T}_{13} + (b_{13})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(13)(j)}T_{13}^{*}\mathbb{G}_{j}\right)$ | 528 |
| $\frac{d\mathbb{T}_{14}}{dt} = -\left((b_{14}^{'})^{(1)} - (r_{14})^{(1)}\right)\mathbb{T}_{14} + (b_{14})^{(1)}\mathbb{T}_{13} + \sum_{j=13}^{15}\left(s_{(14)(j)}T_{14}^{*}\mathbb{G}_{j}\right)$ | 529 |
| $\frac{d\mathbb{T}_{15}}{dt} = -\left((b_{15}^{'})^{(1)} - (r_{15})^{(1)}\right)\mathbb{T}_{15} + (b_{15})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(15)(j)}T_{15}^{*}\mathbb{G}_{j}\right)$ | 530 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}$ Belong to $C^{(2)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable | 531 |
| Proof: Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- | |
| $G_i = G_i^{*} + \mathbb{G}_i \quad , T_i = T_i^{*} + \mathbb{T}_i$ | 532 |
| $\frac{\partial(a_{17}^{''})^{(2)}}{\partial T_{17}}(T_{17}^{*}) = (q_{17})^{(2)}$ , $\frac{\partial(b_i^{''})^{(2)}}{\partial G_j}((G_{19})^{*}) = s_{ij}$ | 533 |
| taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{16}}{dt} = -\left((a_{16}^{'})^{(2)} + (p_{16})^{(2)}\right)\mathbb{G}_{16} + (a_{16})^{(2)}\mathbb{G}_{17} - (q_{16})^{(2)}G_{16}^{*}\mathbb{T}_{17}$ | 534 |
| $\frac{d\mathbb{G}_{17}}{dt} = -\left((a_{17}^{'})^{(2)} + (p_{17})^{(2)}\right)\mathbb{G}_{17} + (a_{17})^{(2)}\mathbb{G}_{16} - (q_{17})^{(2)}G_{17}^{*}\mathbb{T}_{17}$ | 535 |
| $\frac{d\mathbb{G}_{18}}{dt} = -\left((a_{18}^{'})^{(2)} + (p_{18})^{(2)}\right)\mathbb{G}_{18} + (a_{18})^{(2)}\mathbb{G}_{17} - (q_{18})^{(2)}G_{18}^{*}\mathbb{T}_{17}$ | 536 |
| $\frac{d\mathbb{T}_{16}}{dt} = -\left((b_{16}^{'})^{(2)} - (r_{16})^{(2)}\right)\mathbb{T}_{16} + (b_{16})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(16)(j)}T_{16}^{*}\mathbb{G}_{j}\right)$ | 537 |
| $\frac{d\mathbb{T}_{17}}{dt} = -\left((b_{17}^{'})^{(2)} - (r_{17})^{(2)}\right)\mathbb{T}_{17} + (b_{17})^{(2)}\mathbb{T}_{16} + \sum_{j=16}^{18}\left(s_{(17)(j)}T_{17}^{*}\mathbb{G}_{j}\right)$ | 538 |
| $\frac{d\mathbb{T}_{18}}{dt} = -\left((b_{18}^{'})^{(2)} - (r_{18})^{(2)}\right)\mathbb{T}_{18} + (b_{18})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(18)(j)}T_{18}^{*}\mathbb{G}_{j}\right)$ | 539 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(3)}$ *and* $(b_i^{''})^{(3)}$ Belong to $C^{(3)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- | 540 |

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| $G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br>$\frac{\partial (a_{21}^{''})^{(3)}}{\partial T_{21}}(T_{21}^*) = (q_{21})^{(3)} \ , \ \frac{\partial (b_i^{''})^{(3)}}{\partial G_j}((G_{23})^*) = s_{ij}$ | |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{20}}{dt} = -\left((a_{20}^{'})^{(3)} + (p_{20})^{(3)}\right)\mathbb{G}_{20} + (a_{20})^{(3)}\mathbb{G}_{21} - (q_{20})^{(3)}G_{20}^*\mathbb{T}_{21}$ | 541 |
| $\frac{d\mathbb{G}_{21}}{dt} = -\left((a_{21}^{'})^{(3)} + (p_{21})^{(3)}\right)\mathbb{G}_{21} + (a_{21})^{(3)}\mathbb{G}_{20} - (q_{21})^{(3)}G_{21}^*\mathbb{T}_{21}$ | 542 |
| $\frac{d\mathbb{G}_{22}}{dt} = -\left((a_{22}^{'})^{(3)} + (p_{22})^{(3)}\right)\mathbb{G}_{22} + (a_{22})^{(3)}\mathbb{G}_{21} - (q_{22})^{(3)}G_{22}^*\mathbb{T}_{21}$ | 543 |
| $\frac{d\mathbb{T}_{20}}{dt} = -\left((b_{20}^{'})^{(3)} - (r_{20})^{(3)}\right)\mathbb{T}_{20} + (b_{20})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(20)(j)}T_{20}^*\mathbb{G}_j\right)$ | 544 |
| $\frac{d\mathbb{T}_{21}}{dt} = -\left((b_{21}^{'})^{(3)} - (r_{21})^{(3)}\right)\mathbb{T}_{21} + (b_{21})^{(3)}\mathbb{T}_{20} + \sum_{j=20}^{22}\left(s_{(21)(j)}T_{21}^*\mathbb{G}_j\right)$ | 545 |
| $\frac{d\mathbb{T}_{22}}{dt} = -\left((b_{22}^{'})^{(3)} - (r_{22})^{(3)}\right)\mathbb{T}_{22} + (b_{22})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(22)(j)}T_{22}^*\mathbb{G}_j\right)$ | 546 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(4)}$ *and* $(b_i^{''})^{(4)}$ Belong to $C^{(4)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 547 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br>$\frac{\partial (a_{25}^{''})^{(4)}}{\partial T_{25}}(T_{25}^*) = (q_{25})^{(4)} \ , \ \frac{\partial (b_i^{''})^{(4)}}{\partial G_j}((G_{27})^*) = s_{ij}$ | 548 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{24}}{dt} = -\left((a_{24}^{'})^{(4)} + (p_{24})^{(4)}\right)\mathbb{G}_{24} + (a_{24})^{(4)}\mathbb{G}_{25} - (q_{24})^{(4)}G_{24}^*\mathbb{T}_{25}$ | 549 |
| $\frac{d\mathbb{G}_{25}}{dt} = -\left((a_{25}^{'})^{(4)} + (p_{25})^{(4)}\right)\mathbb{G}_{25} + (a_{25})^{(4)}\mathbb{G}_{24} - (q_{25})^{(4)}G_{25}^*\mathbb{T}_{25}$ | 550 |
| $\frac{d\mathbb{G}_{26}}{dt} = -\left((a_{26}^{'})^{(4)} + (p_{26})^{(4)}\right)\mathbb{G}_{26} + (a_{26})^{(4)}\mathbb{G}_{25} - (q_{26})^{(4)}G_{26}^*\mathbb{T}_{25}$ | 551 |
| $\frac{d\mathbb{T}_{24}}{dt} = -\left((b_{24}^{'})^{(4)} - (r_{24})^{(4)}\right)\mathbb{T}_{24} + (b_{24})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(24)(j)}T_{24}^*\mathbb{G}_j\right)$ | 552 |
| $\frac{d\mathbb{T}_{25}}{dt} = -\left((b_{25}^{'})^{(4)} - (r_{25})^{(4)}\right)\mathbb{T}_{25} + (b_{25})^{(4)}\mathbb{T}_{24} + \sum_{j=24}^{26}\left(s_{(25)(j)}T_{25}^*\mathbb{G}_j\right)$ | 553 |
| $\frac{d\mathbb{T}_{26}}{dt} = -\left((b_{26}^{'})^{(4)} - (r_{26})^{(4)}\right)\mathbb{T}_{26} + (b_{26})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(26)(j)}T_{26}^*\mathbb{G}_j\right)$ | 554 |

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| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 5:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(5)}$ *and* $(b_i^{''})^{(5)}$ Belong to $C^{(5)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 555 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial (a_{29}^{''})^{(5)}}{\partial T_{29}}(T_{29}^*) = (q_{29})^{(5)}$ , $\frac{\partial (b_i^{''})^{(5)}}{\partial G_j}((G_{31})^*) = s_{ij}$ | 556 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{28}}{dt} = -\left((a_{28}^{'})^{(5)} + (p_{28})^{(5)}\right)\mathbb{G}_{28} + (a_{28})^{(5)}\mathbb{G}_{29} - (q_{28})^{(5)}G_{28}^*\mathbb{T}_{29}$ | 557 |
| $\frac{d\mathbb{G}_{29}}{dt} = -\left((a_{29}^{'})^{(5)} + (p_{29})^{(5)}\right)\mathbb{G}_{29} + (a_{29})^{(5)}\mathbb{G}_{28} - (q_{29})^{(5)}G_{29}^*\mathbb{T}_{29}$ | 558 |
| $\frac{d\mathbb{G}_{30}}{dt} = -\left((a_{30}^{'})^{(5)} + (p_{30})^{(5)}\right)\mathbb{G}_{30} + (a_{30})^{(5)}\mathbb{G}_{29} - (q_{30})^{(5)}G_{30}^*\mathbb{T}_{29}$ | 559 |
| $\frac{d\mathbb{T}_{28}}{dt} = -\left((b_{28}^{'})^{(5)} - (r_{28})^{(5)}\right)\mathbb{T}_{28} + (b_{28})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(28)(j)}T_{28}^*\mathbb{G}_j\right)$ | 560 |
| $\frac{d\mathbb{T}_{29}}{dt} = -\left((b_{29}^{'})^{(5)} - (r_{29})^{(5)}\right)\mathbb{T}_{29} + (b_{29})^{(5)}\mathbb{T}_{28} + \sum_{j=28}^{30}\left(s_{(29)(j)}T_{29}^*\mathbb{G}_j\right)$ | 561 |
| $\frac{d\mathbb{T}_{30}}{dt} = -\left((b_{30}^{'})^{(5)} - (r_{30})^{(5)}\right)\mathbb{T}_{30} + (b_{30})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(30)(j)}T_{30}^*\mathbb{G}_j\right)$ | 562 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 6:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(6)}$ *and* $(b_i^{''})^{(6)}$ Belong to $C^{(6)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 563 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial (a_{33}^{''})^{(6)}}{\partial T_{33}}(T_{33}^*) = (q_{33})^{(6)}$ , $\frac{\partial (b_i^{''})^{(6)}}{\partial G_j}((G_{35})^*) = s_{ij}$ | 564 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{32}}{dt} = -\left((a_{32}^{'})^{(6)} + (p_{32})^{(6)}\right)\mathbb{G}_{32} + (a_{32})^{(6)}\mathbb{G}_{33} - (q_{32})^{(6)}G_{32}^*\mathbb{T}_{33}$ | 565 |
| $\frac{d\mathbb{G}_{33}}{dt} = -\left((a_{33}^{'})^{(6)} + (p_{33})^{(6)}\right)\mathbb{G}_{33} + (a_{33})^{(6)}\mathbb{G}_{32} - (q_{33})^{(6)}G_{33}^*\mathbb{T}_{33}$ | 566 |

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| $\frac{d\mathbb{G}_{34}}{dt} = -\left((a'_{34})^{(6)} + (p_{34})^{(6)}\right)\mathbb{G}_{34} + (a_{34})^{(6)}\mathbb{G}_{33} - (q_{34})^{(6)}G^*_{34}\mathbb{T}_{33}$ | 567 |
| $\frac{d\mathbb{T}_{32}}{dt} = -\left((b'_{32})^{(6)} - (r_{32})^{(6)}\right)\mathbb{T}_{32} + (b_{32})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(32)(j)}T^*_{32}\mathbb{G}_j\right)$ | 568 |
| $\frac{d\mathbb{T}_{33}}{dt} = -\left((b'_{33})^{(6)} - (r_{33})^{(6)}\right)\mathbb{T}_{33} + (b_{33})^{(6)}\mathbb{T}_{32} + \sum_{j=32}^{34}\left(s_{(33)(j)}T^*_{33}\mathbb{G}_j\right)$ | 569 |
| $\frac{d\mathbb{T}_{34}}{dt} = -\left((b'_{34})^{(6)} - (r_{34})^{(6)}\right)\mathbb{T}_{34} + (b_{34})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(34)(j)}T^*_{34}\mathbb{G}_j\right)$ | 570 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 7:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(7)}$ *and* $(b_i^{''})^{(7)}$ Belong to $C^{(7)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 571 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial(a''_{37})^{(7)}}{\partial T_{37}}(T^*_{37}) = (q_{37})^{(7)}$ , $\frac{\partial(b_i^{''})^{(7)}}{\partial G_j}\left((G_{39})^{**}\right) = s_{ij}$ | 572 |
| Then taking into account equations and neglecting the terms of power 2, we obtain from | |
| $\frac{d\mathbb{G}_{36}}{dt} = -\left((a'_{36})^{(7)} + (p_{36})^{(7)}\right)\mathbb{G}_{36} + (a_{36})^{(7)}\mathbb{G}_{37} - (q_{36})^{(7)}G^*_{36}\mathbb{T}_{37}$ | 573 |
| $\frac{d\mathbb{G}_{37}}{dt} = -\left((a'_{37})^{(7)} + (p_{37})^{(7)}\right)\mathbb{G}_{37} + (a_{37})^{(7)}\mathbb{G}_{36} - (q_{37})^{(7)}G^*_{37}\mathbb{T}_{37}$ | 574 |
| $\frac{d\mathbb{G}_{38}}{dt} = -\left((a'_{38})^{(7)} + (p_{38})^{(7)}\right)\mathbb{G}_{38} + (a_{38})^{(7)}\mathbb{G}_{37} - (q_{38})^{(7)}G^*_{38}\mathbb{T}_{37}$ | 575 |
| $\frac{d\mathbb{T}_{36}}{dt} = -\left((b'_{36})^{(7)} - (r_{36})^{(7)}\right)\mathbb{T}_{36} + (b_{36})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(36)(j)}T^*_{36}\mathbb{G}_j\right)$ | 576 |
| $\frac{d\mathbb{T}_{37}}{dt} = -\left((b'_{37})^{(7)} - (r_{37})^{(7)}\right)\mathbb{T}_{37} + (b_{37})^{(7)}\mathbb{T}_{36} + \sum_{j=36}^{38}\left(s_{(37)(j)}T^*_{37}\mathbb{G}_j\right)$ | 578 |
| $\frac{d\mathbb{T}_{38}}{dt} = -\left((b'_{38})^{(7)} - (r_{38})^{(7)}\right)\mathbb{T}_{38} + (b_{38})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(38)(j)}T^*_{38}\mathbb{G}_j\right)$ | 579 |
| Obviously, these values represent an equilibrium solution<br><br>**ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 8:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(8)}$ *and* $(b_i^{''})^{(8)}$ Belong to $C^{(8)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable. | |

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| Proof: Denote | |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- $G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$ $\frac{\partial (a_{41}^{''})^{(8)}}{\partial T_{41}}(T_{41}^*) = (q_{41})^{(8)} \ , \frac{\partial (b_i^{''})^{(8)}}{\partial G_j}((G_{43})^*) = s_{ij}$ | 580 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{40}}{dt} = -\left((a_{40}^{'})^{(8)} + (p_{40})^{(8)}\right)\mathbb{G}_{40} + (a_{40})^{(8)}\mathbb{G}_{41} - (q_{40})^{(8)}G_{40}^*\mathbb{T}_{41}$ | 581 |
| $\frac{d\mathbb{G}_{41}}{dt} = -\left((a_{41}^{'})^{(8)} + (p_{41})^{(8)}\right)\mathbb{G}_{41} + (a_{41})^{(8)}\mathbb{G}_{40} - (q_{41})^{(8)}G_{41}^*\mathbb{T}_{41}$ | 582 |
| $\frac{d\mathbb{G}_{42}}{dt} = -\left((a_{42}^{'})^{(8)} + (p_{42})^{(8)}\right)\mathbb{G}_{42} + (a_{42})^{(8)}\mathbb{G}_{41} - (q_{42})^{(8)}G_{42}^*\mathbb{T}_{41}$ | 583 |
| $\frac{d\mathbb{T}_{40}}{dt} = -\left((b_{40}^{'})^{(8)} - (r_{40})^{(8)}\right)\mathbb{T}_{40} + (b_{40})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(40)(j)}T_{40}^*\mathbb{G}_j\right)$ | 584 |
| $\frac{d\mathbb{T}_{41}}{dt} = -\left((b_{41}^{'})^{(8)} - (r_{41})^{(8)}\right)\mathbb{T}_{41} + (b_{41})^{(8)}\mathbb{T}_{40} + \sum_{j=40}^{42}\left(s_{(41)(j)}T_{41}^*\mathbb{G}_j\right)$ | 585 |
| $\frac{d\mathbb{T}_{42}}{dt} = -\left((b_{42}^{'})^{(8)} - (r_{42})^{(8)}\right)\mathbb{T}_{42} + (b_{42})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(42)(j)}T_{42}^*\mathbb{G}_j\right)$ | 586 |
| **ASYMPTOTIC STABILITY ANALYSIS** **Theorem 9:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(9)}$ *and* $(b_i^{''})^{(9)}$ Belong to $C^{(9)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable. **Proof: Denote** | 586 A |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- $G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$ $\frac{\partial (a_{45}^{''})^{(9)}}{\partial T_{45}}(T_{45}^*) = (q_{45})^{(9)} \ , \frac{\partial (b_i^{''})^{(9)}}{\partial G_j}((G_{47})^*) = s_{ij}$ | |
| Then taking into account equations 89 to 99 and neglecting the terms of power 2, we obtain from 99 to 44 | |
| $\frac{d\mathbb{G}_{44}}{dt} = -\left((a_{44}^{'})^{(9)} + (p_{44})^{(9)}\right)\mathbb{G}_{44} + (a_{44})^{(9)}\mathbb{G}_{45} - (q_{44})^{(9)}G_{44}^*\mathbb{T}_{45}$ | 586 B |
| $\frac{d\mathbb{G}_{45}}{dt} = -\left((a_{45}^{'})^{(9)} + (p_{45})^{(9)}\right)\mathbb{G}_{45} + (a_{45})^{(9)}\mathbb{G}_{44} - (q_{45})^{(9)}G_{45}^*\mathbb{T}_{45}$ | 586 C |
| $\frac{d\mathbb{G}_{46}}{dt} = -\left((a_{46}^{'})^{(9)} + (p_{46})^{(9)}\right)\mathbb{G}_{46} + (a_{46})^{(9)}\mathbb{G}_{45} - (q_{46})^{(9)}G_{46}^*\mathbb{T}_{45}$ | 586 D |

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|---|---|
| $\frac{d\mathbb{T}_{44}}{dt} = -\left((b'_{44})^{(9)} - (r_{44})^{(9)}\right)\mathbb{T}_{44} + (b_{44})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(44)(j)}T^*_{44}\mathbb{G}_j\right)$ | 586 E |
| $\frac{d\mathbb{T}_{45}}{dt} = -\left((b'_{45})^{(9)} - (r_{45})^{(9)}\right)\mathbb{T}_{45} + (b_{45})^{(9)}\mathbb{T}_{44} + \sum_{j=44}^{46}\left(s_{(45)(j)}T^*_{45}\mathbb{G}_j\right)$ | 586 F |
| $\frac{d\mathbb{T}_{46}}{dt} = -\left((b'_{46})^{(9)} - (r_{46})^{(9)}\right)\mathbb{T}_{46} + (b_{46})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(46)(j)}T^*_{46}\mathbb{G}_j\right)$ | 586 G |
| **The characteristic equation of this system is** | 587 |
| $\left((\lambda)^{(1)} + (b'_{15})^{(1)} - (r_{15})^{(1)}\right)\{\left((\lambda)^{(1)} + (a'_{15})^{(1)} + (p_{15})^{(1)}\right)$ $\left[\left(\left((\lambda)^{(1)} + (a'_{13})^{(1)} + (p_{13})^{(1)}\right)(q_{14})^{(1)}G^*_{14} + (a_{14})^{(1)}(q_{13})^{(1)}G^*_{13}\right)\right]$ $\left(\left((\lambda)^{(1)} + (b'_{13})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(14)}T^*_{14} + (b_{14})^{(1)}s_{(13),(14)}T^*_{14}\right)$ $+\left(\left((\lambda)^{(1)} + (a'_{14})^{(1)} + (p_{14})^{(1)}\right)(q_{13})^{(1)}G^*_{13} + (a_{13})^{(1)}(q_{14})^{(1)}G^*_{14}\right)$ $\left(\left((\lambda)^{(1)} + (b'_{13})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(13)}T^*_{14} + (b_{14})^{(1)}s_{(13),(13)}T^*_{13}\right)$ $\left(\left((\lambda)^{(1)}\right)^2 + \left((a'_{13})^{(1)} + (a'_{14})^{(1)} + (p_{13})^{(1)} + (p_{14})^{(1)}\right)(\lambda)^{(1)}\right)$ $\left(\left((\lambda)^{(1)}\right)^2 + \left((b'_{13})^{(1)} + (b'_{14})^{(1)} - (r_{13})^{(1)} + (r_{14})^{(1)}\right)(\lambda)^{(1)}\right)$ $+\left(\left((\lambda)^{(1)}\right)^2 + \left((a'_{13})^{(1)} + (a'_{14})^{(1)} + (p_{13})^{(1)} + (p_{14})^{(1)}\right)(\lambda)^{(1)}\right)(q_{15})^{(1)}G_{15}$ $+\left((\lambda)^{(1)} + (a'_{13})^{(1)} + (p_{13})^{(1)}\right)\left((a_{15})^{(1)}(q_{14})^{(1)}G^*_{14} + (a_{14})^{(1)}(a_{15})^{(1)}(q_{13})^{(1)}G^*_{13}\right)$ $\left(\left((\lambda)^{(1)} + (b'_{13})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(15)}T^*_{14} + (b_{14})^{(1)}s_{(13),(15)}T^*_{13}\right)\} = 0$ $+$ $\left((\lambda)^{(2)} + (b'_{18})^{(2)} - (r_{18})^{(2)}\right)\{\left((\lambda)^{(2)} + (a'_{18})^{(2)} + (p_{18})^{(2)}\right)$ $\left[\left(\left((\lambda)^{(2)} + (a'_{16})^{(2)} + (p_{16})^{(2)}\right)(q_{17})^{(2)}\mathrm{G}^*_{17} + (a_{17})^{(2)}(q_{16})^{(2)}\mathrm{G}^*_{16}\right)\right]$ $\left(\left((\lambda)^{(2)} + (b'_{16})^{(2)} - (r_{16})^{(2)}\right)s_{(17),(17)}\mathrm{T}^*_{17} + (b_{17})^{(2)}s_{(16),(17)}\mathrm{T}^*_{17}\right)$ $+\left(\left((\lambda)^{(2)} + (a'_{17})^{(2)} + (p_{17})^{(2)}\right)(q_{16})^{(2)}\mathrm{G}^*_{16} + (a_{16})^{(2)}(q_{17})^{(2)}\mathrm{G}^*_{17}\right)$ $\left(\left((\lambda)^{(2)} + (b'_{16})^{(2)} - (r_{16})^{(2)}\right)s_{(17),(16)}\mathrm{T}^*_{17} + (b_{17})^{(2)}s_{(16),(16)}\mathrm{T}^*_{16}\right)$ $\left(\left((\lambda)^{(2)}\right)^2 + \left((a'_{16})^{(2)} + (a'_{17})^{(2)} + (p_{16})^{(2)} + (p_{17})^{(2)}\right)(\lambda)^{(2)}\right)$ $\left(\left((\lambda)^{(2)}\right)^2 + \left((b'_{16})^{(2)} + (b'_{17})^{(2)} - (r_{16})^{(2)} + (r_{17})^{(2)}\right)(\lambda)^{(2)}\right)$ | |

$$+\left(\left((\lambda)^{(2)}\right)^2+\left((a_{16}')^{(2)}+(a_{17}')^{(2)}+(p_{16})^{(2)}+(p_{17})^{(2)}\right)(\lambda)^{(2)}\right)(q_{18})^{(2)}\mathrm{G}_{18}$$

$$+\left((\lambda)^{(2)}+(a_{16}')^{(2)}+(p_{16})^{(2)}\right)\left((a_{18})^{(2)}(q_{17})^{(2)}\mathrm{G}_{17}^{*}+(a_{17})^{(2)}(a_{18})^{(2)}(q_{16})^{(2)}\mathrm{G}_{16}^{*}\right)$$

$$\left(\left((\lambda)^{(2)}+(b_{16}')^{(2)}-(r_{16})^{(2)}\right)s_{(17),(18)}\mathrm{T}_{17}^{*}+(b_{17})^{(2)}s_{(16),(18)}\mathrm{T}_{16}^{*}\right)\}=0$$

$$+$$

$$\left((\lambda)^{(3)}+(b_{22}')^{(3)}-(r_{22})^{(3)}\right)\{\left((\lambda)^{(3)}+(a_{22}')^{(3)}+(p_{22})^{(3)}\right)$$

$$\left[\left(\left((\lambda)^{(3)}+(a_{20}')^{(3)}+(p_{20})^{(3)}\right)(q_{21})^{(3)}G_{21}^{*}+(a_{21})^{(3)}(q_{20})^{(3)}G_{20}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(21)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(21)}T_{21}^{*}\right)$$

$$+\left(\left((\lambda)^{(3)}+(a_{21}')^{(3)}+(p_{21})^{(3)}\right)(q_{20})^{(3)}G_{20}^{*}+(a_{20})^{(3)}(q_{21})^{(1)}G_{21}^{*}\right)$$

$$\left(\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(20)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(20)}T_{20}^{*}\right)$$

$$\left(\left((\lambda)^{(3)}\right)^2+\left((a_{20}')^{(3)}+(a_{21}')^{(3)}+(p_{20})^{(3)}+(p_{21})^{(3)}\right)(\lambda)^{(3)}\right)$$

$$\left(\left((\lambda)^{(3)}\right)^2+\left((b_{20}')^{(3)}+(b_{21}')^{(3)}-(r_{20})^{(3)}+(r_{21})^{(3)}\right)(\lambda)^{(3)}\right)$$

$$+\left(\left((\lambda)^{(3)}\right)^2+\left((a_{20}')^{(3)}+(a_{21}')^{(3)}+(p_{20})^{(3)}+(p_{21})^{(3)}\right)(\lambda)^{(3)}\right)(q_{22})^{(3)}G_{22}$$

$$+\left((\lambda)^{(3)}+(a_{20}')^{(3)}+(p_{20})^{(3)}\right)\left((a_{22})^{(3)}(q_{21})^{(3)}G_{21}^{*}+(a_{21})^{(3)}(a_{22})^{(3)}(q_{20})^{(3)}G_{20}^{*}\right)$$

$$\left(\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(22)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(22)}T_{20}^{*}\right)\}=0$$

$$+$$

$$\left((\lambda)^{(4)}+(b_{26}')^{(4)}-(r_{26})^{(4)}\right)\{\left((\lambda)^{(4)}+(a_{26}')^{(4)}+(p_{26})^{(4)}\right)$$

$$\left[\left(\left((\lambda)^{(4)}+(a_{24}')^{(4)}+(p_{24})^{(4)}\right)(q_{25})^{(4)}G_{25}^{*}+(a_{25})^{(4)}(q_{24})^{(4)}G_{24}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(4)}+(b_{24}')^{(4)}-(r_{24})^{(4)}\right)s_{(25),(25)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(25)}T_{25}^{*}\right)$$

$$+\left(\left((\lambda)^{(4)}+(a_{25}')^{(4)}+(p_{25})^{(4)}\right)(q_{24})^{(4)}G_{24}^{*}+(a_{24})^{(4)}(q_{25})^{(4)}G_{25}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}+(b_{24}')^{(4)}-(r_{24})^{(4)}\right)s_{(25),(24)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(24)}T_{24}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}\right)^2+\left((a_{24}')^{(4)}+(a_{25}')^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)$$

$$\left(\left((\lambda)^{(4)}\right)^2+\left((b_{24}')^{(4)}+(b_{25}')^{(4)}-(r_{24})^{(4)}+(r_{25})^{(4)}\right)(\lambda)^{(4)}\right)$$

$$+\left(\left((\lambda)^{(4)}\right)^2+\left((a_{24}^{'})^{(4)}+(a_{25}^{'})^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)(q_{26})^{(4)}G_{26}$$

$$+\left((\lambda)^{(4)}+(a_{24}^{'})^{(4)}+(p_{24})^{(4)}\right)\left((a_{26})^{(4)}(q_{25})^{(4)}G_{25}^{*}+(a_{25})^{(4)}(a_{26})^{(4)}(q_{24})^{(4)}G_{24}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}+(b_{24}^{'})^{(4)}-(r_{24})^{(4)}\right)s_{(25),(26)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(26)}T_{24}^{*}\right)\}=0$$

$$+$$

$$(\lambda)^{(5)}+(b_{30}^{'})^{(5)}-(r_{30})^{(5)}\}\{\left((\lambda)^{(5)}+(a_{30}^{'})^{(5)}+(p_{30})^{(5)}\right)$$

$$\left[\left(\left((\lambda)^{(5)}+(a_{28}^{'})^{(5)}+(p_{28})^{(5)}\right)(q_{29})^{(5)}G_{29}^{*}+(a_{29})^{(5)}(q_{28})^{(5)}G_{28}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(5)}+(b_{28}^{'})^{(5)}-(r_{28})^{(5)}\right)s_{(29),(29)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(29)}T_{29}^{*}\right)$$

$$+\left(\left((\lambda)^{(5)}+(a_{29}^{'})^{(5)}+(p_{29})^{(5)}\right)(q_{28})^{(5)}G_{28}^{*}+(a_{28})^{(5)}(q_{29})^{(5)}G_{29}^{*}\right)$$

$$\left(\left((\lambda)^{(5)}+(b_{28}^{'})^{(5)}-(r_{28})^{(5)}\right)s_{(29),(28)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(28)}T_{28}^{*}\right)$$

$$\left(\left((\lambda)^{(5)}\right)^2+\left((a_{28}^{'})^{(5)}+(a_{29}^{'})^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)$$

$$\left(\left((\lambda)^{(5)}\right)^2+\left((b_{28}^{'})^{(5)}+(b_{29}^{'})^{(5)}-(r_{28})^{(5)}+(r_{29})^{(5)}\right)(\lambda)^{(5)}\right)$$

$$+\left(\left((\lambda)^{(5)}\right)^2+\left((a_{28}^{'})^{(5)}+(a_{29}^{'})^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)(q_{30})^{(5)}G_{30}$$

$$+\left((\lambda)^{(5)}+(a_{28}^{'})^{(5)}+(p_{28})^{(5)}\right)\left((a_{30})^{(5)}(q_{29})^{(5)}G_{29}^{*}+(a_{29})^{(5)}(a_{30})^{(5)}(q_{28})^{(5)}G_{28}^{*}\right)$$

$$\left(\left((\lambda)^{(5)}+(b_{28}^{'})^{(5)}-(r_{28})^{(5)}\right)s_{(29),(30)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(30)}T_{28}^{*}\right)\}=0$$

$$+$$

$$(\lambda)^{(6)}+(b_{34}^{'})^{(6)}-(r_{34})^{(6)}\}\{\left((\lambda)^{(6)}+(a_{34}^{'})^{(6)}+(p_{34})^{(6)}\right)$$

$$\left[\left(\left((\lambda)^{(6)}+(a_{32}^{'})^{(6)}+(p_{32})^{(6)}\right)(q_{33})^{(6)}G_{33}^{*}+(a_{33})^{(6)}(q_{32})^{(6)}G_{32}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(6)}+(b_{32}^{'})^{(6)}-(r_{32})^{(6)}\right)s_{(33),(33)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(33)}T_{33}^{*}\right)$$

$$+\left(\left((\lambda)^{(6)}+(a_{33}^{'})^{(6)}+(p_{33})^{(6)}\right)(q_{32})^{(6)}G_{32}^{*}+(a_{32})^{(6)}(q_{33})^{(6)}G_{33}^{*}\right)$$

$$\left(\left((\lambda)^{(6)}+(b_{32}^{'})^{(6)}-(r_{32})^{(6)}\right)s_{(33),(32)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(32)}T_{32}^{*}\right)$$

$$\left(\left((\lambda)^{(6)}\right)^2+\left((a_{32}^{'})^{(6)}+(a_{33}^{'})^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\right)$$

$$\left(\left((\lambda)^{(6)}\right)^2+\left((b_{32}^{'})^{(6)}+(b_{33}^{'})^{(6)}-(r_{32})^{(6)}+(r_{33})^{(6)}\right)(\lambda)^{(6)}\right)$$

$$+\left(\left((\lambda)^{(6)}\right)^2+\left((a'_{32})^{(6)}+(a'_{33})^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\right)(q_{34})^{(6)}G_{34}$$

$$+\left((\lambda)^{(6)}+(a'_{32})^{(6)}+(p_{32})^{(6)}\right)\left((a_{34})^{(6)}(q_{33})^{(6)}G^*_{33}+(a_{33})^{(6)}(a_{34})^{(6)}(q_{32})^{(6)}G^*_{32}\right)$$

$$\left((\lambda)^{(6)}+(b'_{32})^{(6)}-(r_{32})^{(6)}\right)s_{(33),(34)}T^*_{33}+(b_{33})^{(6)}s_{(32),(34)}T^*_{32}\Big)\}=0$$

$$+$$

$$(\lambda)^{(7)}+(b'_{38})^{(7)}-(r_{38})^{(7)}\}\{\left((\lambda)^{(7)}+(a'_{38})^{(7)}+(p_{38})^{(7)}\right)$$

$$\left[\left(\left((\lambda)^{(7)}+(a'_{36})^{(7)}+(p_{36})^{(7)}\right)(q_{37})^{(7)}G^*_{37}+(a_{37})^{(7)}(q_{36})^{(7)}G^*_{36}\right)\right]$$

$$\left(\left((\lambda)^{(7)}+(b'_{36})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(37)}T^*_{37}+(b_{37})^{(7)}s_{(36),(37)}T^*_{37}\right)$$

$$+\left(\left((\lambda)^{(7)}+(a'_{37})^{(7)}+(p_{37})^{(7)}\right)(q_{36})^{(7)}G^*_{36}+(a_{36})^{(7)}(q_{37})^{(7)}G^*_{37}\right)$$

$$\left(\left((\lambda)^{(7)}+(b'_{36})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(36)}T^*_{37}+(b_{37})^{(7)}s_{(36),(36)}T^*_{36}\right)$$

$$\left(\left((\lambda)^{(7)}\right)^2+\left((a'_{36})^{(7)}+(a'_{37})^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\right)$$

$$\left(\left((\lambda)^{(7)}\right)^2+\left((b'_{36})^{(7)}+(b'_{37})^{(7)}-(r_{36})^{(7)}+(r_{37})^{(7)}\right)(\lambda)^{(7)}\right)$$

$$+\left(\left((\lambda)^{(7)}\right)^2+\left((a'_{36})^{(7)}+(a'_{37})^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\right)(q_{38})^{(7)}G_{38}$$

$$+\left((\lambda)^{(7)}+(a'_{36})^{(7)}+(p_{36})^{(7)}\right)\left((a_{38})^{(7)}(q_{37})^{(7)}G^*_{37}+(a_{37})^{(7)}(a_{38})^{(7)}(q_{36})^{(7)}G^*_{36}\right)$$

$$\left(\left((\lambda)^{(7)}+(b'_{36})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(38)}T^*_{37}+(b_{37})^{(7)}s_{(36),(38)}T^*_{36}\right)\}=0$$

$$+$$

$$(\lambda)^{(8)}+(b'_{42})^{(8)}-(r_{42})^{(8)}\}\{\left((\lambda)^{(8)}+(a'_{42})^{(8)}+(p_{42})^{(8)}\right)$$

$$\left[\left(\left((\lambda)^{(8)}+(a'_{40})^{(8)}+(p_{40})^{(8)}\right)(q_{41})^{(8)}G^*_{41}+(a_{41})^{(8)}(q_{40})^{(8)}G^*_{40}\right)\right]$$

$$\left(\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(41)}T^*_{41}+(b_{41})^{(8)}s_{(40),(41)}T^*_{41}\right)$$

$$+\left(\left((\lambda)^{(8)}+(a'_{41})^{(8)}+(p_{41})^{(8)}\right)(q_{40})^{(8)}G^*_{40}+(a_{40})^{(8)}(q_{41})^{(8)}G^*_{41}\right)$$

$$\left(\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(40)}T^*_{41}+(b_{41})^{(8)}s_{(40),(40)}T^*_{40}\right)$$

$$\left(\left((\lambda)^{(8)}\right)^2+\left((a'_{40})^{(8)}+(a'_{41})^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\right)$$

$$\left((\lambda)^{(8)}\right)^2+\left((b'_{40})^{(8)}+(b'_{41})^{(8)}-(r_{40})^{(8)}+(r_{41})^{(8)}\right)(\lambda)^{(8)}\Big)$$

$$+\left(\left((\lambda)^{(8)}\right)^2+\left((a'_{40})^{(8)}+(a'_{41})^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\right)(q_{42})^{(8)}G_{42}$$

$$+\left((\lambda)^{(8)}+(a'_{40})^{(8)}+(p_{40})^{(8)}\right)\left((a_{42})^{(8)}(q_{41})^{(8)}G^*_{41}+(a_{41})^{(8)}(a_{42})^{(8)}(q_{40})^{(8)}G^*_{40}\right)$$

$$\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(42)}T^*_{41}+(b_{41})^{(8)}s_{(40),(42)}T^*_{40}\Big)\}=0$$

$$+$$

$$(\lambda)^{(9)}+(b'_{46})^{(9)}-(r_{46})^{(9)}\}\{\left((\lambda)^{(9)}+(a'_{46})^{(9)}+(p_{46})^{(9)}\right)$$

$$\left[\left(\left((\lambda)^{(9)}+(a'_{44})^{(9)}+(p_{44})^{(9)}\right)(q_{45})^{(9)}G^*_{45}+(a_{45})^{(9)}(q_{44})^{(9)}G^*_{44}\right)\right]$$

$$\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(45)}T^*_{45}+(b_{45})^{(9)}s_{(44),(45)}T^*_{45}\Big)$$

$$+\left(\left((\lambda)^{(9)}+(a'_{45})^{(9)}+(p_{45})^{(9)}\right)(q_{44})^{(9)}G^*_{44}+(a_{44})^{(9)}(q_{45})^{(9)}G^*_{45}\right)$$

$$\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(44)}T^*_{45}+(b_{45})^{(9)}s_{(44),(44)}T^*_{44}\Big)$$

$$\left((\lambda)^{(9)}\right)^2+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\Big)$$

$$\left((\lambda)^{(9)}\right)^2+\left((b'_{44})^{(9)}+(b'_{45})^{(9)}-(r_{44})^{(9)}+(r_{45})^{(9)}\right)(\lambda)^{(9)}\Big)$$

$$+\left(\left((\lambda)^{(9)}\right)^2+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\right)(q_{46})^{(9)}G_{46}$$

$$+\left((\lambda)^{(9)}+(a'_{44})^{(9)}+(p_{44})^{(9)}\right)\left((a_{46})^{(9)}(q_{45})^{(9)}G^*_{45}+(a_{45})^{(9)}(a_{46})^{(9)}(q_{44})^{(9)}G^*_{44}\right)$$

$$\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(46)}T^*_{45}+(b_{45})^{(9)}s_{(44),(46)}T^*_{44}\Big)\}=0$$

**And as one sees, all the coefficients are positive. It follows that all the roots have negative real part, and this proves the theorem.**

Note: **Same dovetailing explanation holds good for d/dt (partial differential with respect to t), d/dt,d^2/dt^2 (acceralation: double dot). Such an exposition is helpful in optimisation problems, duality of motion, wave motion a homogeneous functions studies, and control theory which we intend to incorporate in future.** It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the

concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

# Persistence, Coexistence And Thrusty Permanence For Interacting Structured Particles, Transactions In Finance, Econometric Productivity Functions, Environment, Interfaces, Cosmological Interactions Et Al

## Josef Hofbauer And Sebastian J. Schreiber Model Thereof: A Functional Appraisal

ABSTRACT. **JOSEF HOFBAUER AND SEBASTIAN J. SCHREIBER** studied the dynamics of interacting structured systems (in their case population dynamics) can be modeled by $\frac{dxi}{dt} = A_i\ (x)\ xi$ where $x_i\ \epsilon\ R^n_{i,}\ x = (x_1, \ldots, x_k)$, and $A_i(x)$ are matrices with non-negative off- diagonal entries. These models are permanent if there exists a positive global attractor and are robustly permanent if they remain permanent following perturbations of $A_i\ (x)$. Necessary and sufficient conditions for robust permanence are derived using dominant Lyapunov exponents $X_i\ (\mu)$ of the $A_i\ (x)$ with respect to invariant measures $\mu$. The necessary conditions requires $\max_i\ X_i(\mu)>0$ for all ergodic measures with support in the boundary of the non-negative cone. The sufficient condition requires that the boundary admits a Morse decomposition such that $\max_i$ Xi$(\mu)>0$ for all invariant measures $\mu$ supported by a component of the Morse decomposition. When the Morse components are Axiom A, uniquely ergodic, or support all but one population, the necessary and sufficient conditions are equivalent. Applications to spatial ecology, epidemiology, and gene networks are given. Appeared in *Journal of Differential Equations*, 248, 1955-1971 (2010) It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## INTRODUCTION

A fundamental issue in population biology is what are the minimal conditions to ensure the long-term survivorship for interacting populations whether they be viral particles, bio-chemicals, plants, or animals. When these conditions are met the interacting populations are said to persist or coexist. Since the pioneering work of Lotka and Volterra on competitive and predator prey interactions, Thompson, Nicholson, and Bailey on host-parasite interactions, and

Kermack and McKendrick on disease outbreaks, nonlinear difference and differential equations have been used to understand conditions for population persistence [1-5]. One particularly important form of persistence for deterministic models is *permanence* or *uniform persistence* which corresponds to the existence of a global attractor bounded away from extinction of one or more species. When such an attractor exists, interacting populations are able to recover from "vigorous shake ups" of the of the population state. Permanence has been characterized from a topological perspective [6-8] and with average Lyapunov functions [9-12] for differential equation models of the form

$$(1) \qquad \frac{dxi}{dt} = x_i f_i(x) \qquad i = 1, \ldots, k$$

Where $x_i$ is the abundance of population $i$, $x = (x_1, \ldots, x_k)$, and $f_i(x)$ is the per-capita growth of population $i$.

An important extension of the concept of the permanence is *robust permanence* [13-16] which requires that (1) remain permanent after sufficiently small perturbations of the per-capita growth rates $f_i$. The importance of robust permanence as one crude form of structural stability for population models are approximations to reality. Consequently, if nearby models (e.g. more realistic models) are not permanent despite the focal model being permanent, and then one can draw few conclusions about the persistence of the biological system being approximated by the model. One can view robust permanence as one crude form of structural stability for debits-credits particles, propositions, transactions in the celestial subterranean realm and ceratoid dualism, finance like debiting say a current account and crediting to another account, or recrediting to the account which had been earlier debited, further conduction of the cash credit account, black hole formation, guzzling of inventories for the output, creation of income from output, reinvestment of income to investment,gormadizing of matter by black holes, and emission of radiation therefore, gamma ray bursts, in fact most of the transactions and interactions in the universe are conservative and quantum mechanical universe, something like a account acting as predator and prey. More often than not as we have seen from Haimovici model such models could be applied to the problems, propositions, which are not subsumed by the concomitant problems in the bearing. the population models. For this perspective, it is not unexpected that there are permanent systems that cannot be approximated by robustly permanent systems [17]. In [14, 16], criteria for robust permanence were developed with respect to the average per-capita growth rates $\int f_i(x)d\mu(x)$ with respect to invariant probability measures $\mu$ supported on the boundary of the positive orthant (i.e.

where one or more Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates are extinct). Roughly, these criteria for robust permanence require that the average per-capita growth rate (i.e. invasion rate) is positive for some missing species.

While (1) can account for many types of Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates. Interactions, it assumes that all individual within populations are exactly the same. However, theoretical biologists have long recognized that different individuals within a population may be in different states (e.g. different sizes or ages, living in different spatial locations) and these differences can have important consequences for population dynamics [18-24]. To account for how these differences influence persistence of interacting populations, we develop criteria for robust permanence for the dynamics of $k$ interacting structured populations. Roughly, these dynamics correspond to replacing $x_i$ in (1) by vectors and $f_i(x)$ by matrices. In section 2, we describe these models in greater detail and review some basic concepts from dynamical systems theory. In section 3, we introduce the structured analog of average per-capita growth rates and define the concept of robustly unsaturated invariant sets. Necessary and sufficient conditions for an invariant set to be robustly unsaturated are proven. For Axiom A or uniquely ergodic invariant sets, the necessary and sufficient criteria are shown to be equivalent. In section 4, we use Morse decompositions and the criteria for robustly unsaturated invariant sets to develop necessary and sufficient conditions for robust permanence. In sections 5 and 6, we provide several applications of our results to spatially structured ecological models, structured epidemiological models, and models of gene networks. We apply the model to econometrics. To give an example investment needs inventory to be used towards the consummation of output, part of output is to be sold for consolidation of income, and some income might be replenished in the investment and so on. In finance a cash credit account could be debited or credited when there is say purchase of demand bills or discount of usance bills. Cash credit account is debited whenever cheques are issued towards the purchase of raw materials and so on. In fact some of the models have been applied to Organic psychiatry and we have shown how such models are useful in dealing with the propositions of Husserl, Derrida, Deleuze and others.

## 1. MODELS AND ASSUMPTIONS

Let $x_i$ denote the state of the $i$-th population that lies in the non-negative cone $C_i$ of $R^{ni.}$ Define $x=(x_1,\ldots x_k)\epsilon C$ to be the non-negative cone of $R^n$ where $n=\sum_{i=1}^{k} n_i$. Let $C_i^+$ denote positive cone { $x_i \epsilon C_i$: $\prod_i$i noitalupop rof$^j$

$i$ *and* $C^+=\prod_i C^+_i$ , the positive cone for the interacting populations. If $x \mapsto A_i(x)$ is a map into $n_i x n_i$ matrices that describes the growth of population I, then the dynamics of the interacting populations are given by

$$(2)\frac{dx_i}{dt} = A_i(x)x_i \quad i=1,\ldots,k$$

Let $x.t$ denote the solution of (2), with the initial condition x, and more generally the semi flow generated by (2). Before stating our assumption on (2), recall a few definitions from dynamical systems. Given sets $I \subseteq R$ and $K \subseteq C$, let $K.I = \{x.t : t \epsilon I, x \epsilon K\}$. A set $K \subseteq C$ *is invariant* if K.t = K for all t>0, and *forward invariant* if $K.t \subset K$ for all $t > 0$. The *omega limit* set of a set $K \subseteq C$ equals $\omega(K) = \bigcap_{t\geq 0,} \overline{K.[t,\infty)}$. The *alpha limit* set of $K \subseteq C$ is $a(K) = \bigcap_{t\leq 0} \overline{K.(-\infty,t]}$. Given a forward invariant set K, $B \subset K$ is an attractor for the semi flow $x.t$ *restricted* to K provided there exists an open neighborhood $U \subseteq K$ of B such that $\omega(U) = B$. The stable set of a *compact invariant* set $K$ is defined by

$$W^s(K) = \{x \epsilon C : \omega(x) \neq \emptyset \text{ and } \omega(x) \subseteq K\}.$$

The semi flow generated by (2) is *dissipative* if there exists an attractor B with $W^s(B) = C$.

Throughout this paper, we make the following assumptions:

**A1**: $x \mapsto A_i(x)$ is continuous

**A2:** $A_i(x)$ is irreducible and has non-negative off- diagonal entries,

**A3:** $x.t$ is defined for all $t \geq 0$

**A4:** $x.t$ is dissipative with global attractor $\Gamma(A)$.

A1 is a basic regularity assumption. The non-negative of off- diagonal entries in A2 implies that there are no negative feedbacks between individuals of different states in population $i$. This assumption is meet for many types of structured models as discussed in sections 5 and 6. The irreducibility assumption A2 is generally meet and implies that all individuals within a population can pass through all states. In Remark1, we discuss how this irreducibility assumption can be relaxed. A3 ensures that the population dynamics are defined for all future time. A4 requires that

population densities/ abundances eventually are uniformly bounded, a condition that should be meet for an biologically realistic model.

We say that the semi flow of (2) is *permanent* if this semi flow is dissipative and there is a positive attractor $B \subset C^+$ such that $W^s(B) \subset C^+$. We will denote this positive attractor by $\Lambda(A)$. Permanence ensures that populations can recover from rare large perturbations and allows for a diversity of dynamical behaviors.

For K ⊂ C and $\delta > 0$, define the $\delta$- *neighborhood* of K as

$$N\delta(K) = \{x \in C : \backslash x - y < \delta \text{ for some } y \in K\}.$$

We define a $\delta$- *perturbation of (2)* to be a system of the form

$$\frac{dx}{dt} = \tilde{A}(x)x$$

that satisfies assumptions A1-A4, $\Gamma(\tilde{A}) \subset N$ *( $\Gamma(A)$), and* $||A(x) - \tilde{A}(x)|| \leq \delta$ *for all* $x \in N$ *($\Gamma(A)$).* (2) is *robustly permanent* if there exist >0 and $\epsilon > 0$ such that all $\delta$- perturbations of (2) are permanent, and $d(\Lambda(\tilde{A}), C\backslash C^+) > \epsilon$ for all $\delta$- perturbations i.e., there is a common/ uniform region of repulsion around the boundary.

## INUNDATION OF COMPACT SETS

We begin by studying the linear skew product flows on $\Gamma(A) \times R^n_i$ defined by $(x,y).t = (x.t, B_i(t,x)y)$ where $Y(t) = B_i(t,x)$ is the solution to $\frac{dY}{dt}(t) = A_i(x.t)Y(t)$ with y(0) equal to the identify matrix. Our assumption that $A_i$ is irreducible with non- negative off diagonal entries implies that $B_i$ (t,x)$C_i \subset C_i^+$ for all $x$ and t > 0 (see, e.g., [25] ). A result of Ruelle [26, Prop, 3.2] implies that there exist continuous maps $u_i\, v_i : \Gamma(A) \rightarrow C_i^+$ with $u_i(x)| = |v_i(x)| = 1$ such that The line bundle $E_i(x)$ spanned by $u_i(x)$ is invariant i.e. $E_i(x.t) = B_i(x,t)E_i(x)$ for all t ≥ 0.The vector bundle $F_i(x)$ perpendicular to $v_i(x)$ is invariant i.e. $F_i(x.t) = B_i(x,t)F_i(x)$ for all $t \geq 0$.There exist constants $a > 0$ *and* $\beta > 0$ such that(3) $||B_i(t,x)F_i(x)|| \leq a\, exp(-\beta t)||B_i(t,x)E_i(x)||$ For all $x \in K$ *and* $t \geq 0$.**Interaction rates.** Given any x $\epsilon$ $\Gamma$(A), define the *interaction and inundation rate* of *species i at particle/proposition/problem/ state x as* $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right)$ *in* $||B_i(t,x)||$ when $x_i = 0$, $\lambda_i(x)$ provides an upper bound to the rate of growth of population *i* when introduced at infinitesimally small densities. Important properties of this invasion rate are summarized in the following two propositions. For instance, the second proposition implies if $(x.t)_i$ stays bounded away from zero (i.e. population *i* persist), then

$\lambda_i(x)=0$. Alternatively, if $\lambda_i(x)<0$,then population $i$ is doomed to extinction. **Proposition 1.** $\lambda_i(x)$ *satisfies the following properties:* $\lambda_i(x)= limsup_{t\to\infty}\left(\frac{1}{t}\right) In \ |B_i \ (t,x)v|$ *for any* $v>0$ *in* $R^{ni}$, *and* $\lambda_i(x)= limsup_{t\to\infty}\left(\frac{1}{t}\right) \int^{t}_{0}(A_i(x.s)u_i(x.s),u_i(x.s))ds.$

**Proof:** To prove the first property, we first show that

$$\lambda i(x) = \underset{t\to\infty}{limsup}\frac{1}{t} In\ |B_i(t,x)u_i(x)|.$$

To this end, let $v \epsilon R^{ni}$ be any non-zero vector. Since $R^{ni}= E_{i(x)} \ F_i(x)$, we can write $v=au_i(x)+w$ with $a \epsilon R$ and $w \epsilon F_i(x)$. Equation (3) implies

$|B_i(t,x)v| \ \leq \ a|\ B_i(t,x)u_i(x)|+ |B_i(t,x)w| \ \leq \ |\ B_i(t,x)u_i(x)|(a+ \alpha+ exp\ (-\beta t)|\ w|)$. Hence, $\underset{t\to\infty}{limsup}\frac{1}{t} In\ B_i(t,x)\ v| \leq$

$$\underset{t\to\infty}{limsup}\frac{1}{t} In\ B_i\ (t,x)\ u_i\ (x) \leq \lambda_i(x)$$

for all non zero vectors $v \epsilon R^{ni}$. Since $||B_i(t,x)|| = sup|v|=1B_i(t,x)v|$, this inequality implies that $\lambda i(x) = \underset{t\to\infty}{\text{limsup}}\frac{1}{t}$ In $|B_i(t,x)v|$, this inequality implies that $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right)$ In $|t,x)u_i(x)|$. Now let v $\epsilon$ $C_i^{+}$. Then, we can write $v = au_i(x)+ w$ with $a >0$ and $w \epsilon F_i$.

Equation (3) implies

$$|B_i(t,x)v| \ \geq \ a|B_i(t,x)u_i(x)|-|B_i(t,x)w|$$

$$\geq |\ B_i(t,x)u_i(x)|(a- \alpha+ exp\ (-\beta t)|\ w|).$$

Since a>0, this inequality implies that

$$\lambda i(x) \geq \underset{t\to\infty}{\text{limsup}}\frac{1}{t} \text{In}\ |B_i(t,x)v| \ \geq \ \lambda i(x) \geq \underset{t\to\infty}{\text{limsup}}\frac{1}{t} \text{In}\ |B_i(t,x)u_i(x)| = \lambda i(x)$$

To prove the second assertion, let $b_i(t,x) = In\ |B_i(t,x)u_i(x)|$. Invariance of $u_i(x)$ implies that $b_i(t,x)$ is additive:

$$B_i(t+s,x) \ = \ In|B_i(t+s,x)\ u_i(x)|= In\ |B_i(s,x.t)B_i(t,x)u_i(x)$$

$$= \text{In}\ |B_i(s,x.t)\ u_i(x.t)||\ B_i(t,x)\ u_i(x)|$$

$$= \text{In}\ b_i(s,t.x)\ b_i(t,x)$$

Additivity of $b_i(t,x)$ and the fact that $b_i(0,x)$ =0 implies

$$\frac{d}{dt}b_i(t,x) = \lim_{s\to 0}\frac{bi(t+s,x-bi(t,x)}{S-}$$

$$= \lim_{s\to 0}\frac{bi(s,x.t)}{S} = \frac{d}{ds}|_{s=0}\, b_i(s,x.t)$$

=

$$=(A_i(x.t)u_i(x.s),u_i(x.s),$$

The Fundamental Theorem of Calculus implies

$$B_i(t,x)= \int_0^t (A_i(x.s)\, u_i(x.s),\, u_i(x.s))ds\backslash$$

and the second assertion follows.

**Proposition 2:** *For the solutions of (2) we have:* if $x_i>0$, then $\lambda i(x) \le 0$. If $\lambda i(x)<0$, then $\lim_{t\to\infty} (x.t)_i=0$ If $limsup_{t\to\infty}$ $(x.t)_i|>0$, then $\lambda i(x)=0$.

**Proof:** First, assume that x satisfies $x_i>0$. Since the semi flow is dissipative, there exists $\gamma >0$ such that $|x.t| \le \gamma$ for all $t \ge 0$. Proposition 1 and the definition of the skew product flow imply

$$\lambda i(x) = \operatorname{limsup}_{t\to\infty}\frac{1}{t}\,\mathrm{In}\,|B_i(t,x)x_i| = \operatorname{limsup}_{t\to\infty}\frac{1}{t}\,\mathrm{In}\,|\,(x.t)_i\,|$$

$$\le \operatorname{limsup}_{t\to\infty}\frac{In\,\gamma}{t} = 0$$

To prove the second assertion, assume that $\lambda i(x) < 0. If\ x_i=0$, then the invariance of the faces of C imply that $(x.t)_i=0$ for all $t\ge 0$. Alternatively, if $x_i> 0$, then Proposition I and the skew product flow imply

$$\operatorname{limsup}_{t\to\infty}\frac{1}{t}In|(x.t)_i| = \lambda i(x) < 0.$$

The final assertion follows from the first two assertions. 5

**3.2. Invariant measures.** We review some definitions from ergodic theory. Given a Borel probability measure $\mu$ on ***C***, the *support* of $\mu$, denoted supp( $\mu$),is the smallest closed set whose complement has measure zero. A Borel probability measure $\mu$ is called *invariant* for (2) provided that Integral $h(x)d\mu(x) =$ Integral $h(x.t)d\,\mu(x)$for all

$t \geq 0$ and for all bounded continuous functions $h$:C→R. An invariant measure $\mu$ is called *ergodic* provided that $\mu(B)$=0 or 1 for any invariant Borel set *B.*

For an invariant measure $\mu$ define the invasion rate species $i$ with respect to $\mu$ as

$$\lambda_i(\mu) = \int_c \langle \mathrm{A}_i(x)\mathrm{u}_i(x), \mathrm{u}_i(x)\rangle d\mu(x)$$

**Proposition 3:** Let $\mu$ be an *invariant measure.* $\lambda_i(\mu)$ *Satisfies the following properties:* $\lambda_i(\mathrm{x}) = \lim_{t\to\infty} \frac{1}{t}\int_0^t \langle \mathrm{A}_i(x.s)\mathrm{u}_i(x.s), \mathrm{u}_i(x.s)\rangle, ds$ exists for $\mu$ - *almost every* $x$. *Moreover, if* $\mu$ *is ergodic,* $\lambda_i(\mathrm{x}) = \lambda_i(\mu)$ for $\mu$-*almost every* $x$.*If* $\mu$ *is ergodic, then there exists* $I \subset \{1, \dots, k\}$such that $\mu\,(\Pi_{i\in I}\boldsymbol{C}_i^+)$=1 and $\lambda_i(\mu) = 0\ for\ all\ \ i \in I$.

**Proof:** Let $\mu$ be an invariant measure. Define $\mathrm{h}_i(x) = \langle \mathrm{A}_i(x)\mathrm{u}_i(x), \mathrm{u}_i(x)\rangle$. The Birkhoff's' ergodic theorem and Proposition 1 imply

$$\lambda_i(x) = \lim_{t\to\infty} \frac{1}{t}\int_0^t \mathrm{h}_i(x.s)\mathrm{ds}$$

Exists for $\mu$- almost every $x$. Moreover, $\lambda_i(x) = \lambda_i(\mu)$ $\mu$- almost surely if $\mu$ is ergodic.

Assume $\mu$ is ergodic. By ergodicity and invariance of the faces of **C**, there exists $I \subset \{1, \dots, k\}$ such that $\mu\,(\Pi_{i\in I}\boldsymbol{C}_i^+)$=1. Let K $\subset \Pi_{i\in I}\boldsymbol{C}_i^+$ be a compact set such that $\mu(K) > 0$. The Poincare recurrence theorem and the Birkhoff ergodic theorem imply that there is $x \in \Pi_{i\in I}\boldsymbol{C}_i^+$such that $x.\mathrm{t}_n \in K$ for some $\mathrm{t}_n \uparrow \infty$ and $\lambda_i(x) = \lambda_i(\mu)$ for all $i$. The third assertion of Proposition 2 implies $\lambda_i(\mu) = 0$ for all $i \in I$.

**3.3. Robustly unsaturated sets**. Let K $\subset \partial C$ be a compact isolated invariant set for the flow $x.t$ restricted to $\partial C$.*K* is unsaturated if Ws (K) $\subset \partial C$ and K is isolated for $x.t$. If *K* is not unsaturated, then *K* is saturated. *K* is robustly unsaturated for (2) if there exists $\delta > 0$ such that the continuation of K for any $\delta -$ perturbation of (2) is unsaturated.**Theorem 1:** Let *K* be a compact isolated invariant set for $x.t$ restricted to$\partial C$. If one of the following equivalent conditions hold For all invariant measures $\mu$ supported by *K* (4) $\max_{1\leq i\leq k} \lambda_i(\mu) = 0$

There exist $\mathrm{p}_i > 0$ such that (5) $\sum_{1\leq i\leq k} \mathrm{p}_i\,\lambda_i(\mu) = 0$

For all ergodic probability measures supported by *K*.

then $K$ is robustly unsaturated for (2). Alternatively, if $x \to A(x)$ is twice continuously differentiable and $K$ is robustly unsaturated, then (4) holds for all ergodic measures $\mu$ supported by *K*.

**Remark 1**: For some applications (e.g., the disease model considered in section 6.1) it useful to relax the irreducibility assumption A2. For instance, if there exists an open neighborhood U of K such that for each $i$, $A_i(x)$ has a fixed off-diagonal sign pattern for all $x \in U$,then $A_i(x)$ can be decomposed into a finite number, say $m_i$, of irreducible components. For each of these irreducible components, one can define $\lambda_i^j(\mu) = \int_C \langle A_i^j(x) u_i^j(x), u_i^j(x) \rangle \, d\mu(x)$where $A_i^j(x)$ is the submatrix of $A_i(x)$ corresponding to the *j*-th irreducible component of $A_i(x)$ *and* $u_i^j(x)$ is the continuous invariant sub bundle for the irreducible linear cocycle determined by $A_i^j(x)$. If we define $\lambda_i(\mu) = \max_{1 \le j \le m_i} \lambda_i^j(\mu)$, then all of the assertions of Theorem 1 still hold.

**Proof:** To prove the first assertion of the theorem, we need the following lemma. We call an invariant measure $\mu$ for (2) saturated if $\lambda_i(\mu) \le 0$ for all $i$

**Lemma 1**: Let $\delta_n$ be a non –negative sequence that converges to zero as n$\to \infty$. If $\mu_n$ are saturated invariant measures for $\delta_n -$ perturbations of (2), then the weak$^*$ limit points of $\{\mu_n\}_{n=1}^{\infty}$ is a non-empty set consisting of saturated invariant measures for (2).

**Proof:** Let $\delta_n$ be a non-negative sequence that coverages to zero. Let $\mu_n$ be saturated invariant measures for $\delta_n$ perturbations, $\dot{x} = A^n(x)x$, of $\dot{x} = A(x)x$. By weak$^*$ compactness of Borel probability measures supported on $N_1(\Gamma(A))$, there exist weak$^*$ limit points of $\mu_n$ . Let $\mu$ be such a weak$^*$ limit point. Since $\mu_n$ are supported by $N_1(\Gamma(A))$, for all $n$ sufficiently large, $\mu$ is supported by $N_1(\Gamma(A))$.To verify that $\mu$ is an invariant measure for (2),let $h: C \to R$ be a bounded continuous function. Let $t > 0$ and $\in > 0$ be given. Choose $n$ sufficiently large so that $|\int_c h(x)\, d\mu(x) - \int_c h(x) d\mu_n(x)| \le \in$, $|\int_c h(x.t)\, d\mu(x) - \int_c h(x.t) d\mu_n(x)| \le \in$, and $|h(x.t) - h(x^n.t)| \le \in$ for all $x \in N_1(\Gamma(A))$.Then $|\int_c h(x.t) - h(x) d\mu(x)|$ is

$\leq \quad |\int_c h(x.t)\,d\mu\,(x) - \int_c h(x.t)\,d\mu_n\,(x)| + |\int_c h(x.t)\,d\mu_n\,(x) - \int_c h(x)\,d\mu\,(x)|$

$\leq \quad \in + \int_c |h(x.t) - h(x^n.t)|\,d\mu_n\,(x)| + |\int_c h(x^n.t)|\,d\mu_n\,(x) - \int_c h(x)d\mu\,(x)|$

$\leq \quad 2 \in + |\int_c h(x^n.t) - h(x)\,d\mu_n\,(x)| + |\int_c h(x)d\mu_n(x) - \int_c h(x)d\mu\,(x)|$

$\leq \quad 3 \in$

Where the last line follows from the invariance of $\mu_n$ for $x^n.t$. Since $\in > 0$ and t > 0 are arbitrary, we have $\int_c h(x.t)d\mu\,(x) = \int_c h(x)d\mu\,(x)$ for all t >0 and all bounded continuous functions $h: C \to R$. It follows that $\mu$ is an invariant measure for (2). To see that $\mu$ is saturated, define $h_i^n(x) = \langle A_i^n(x)u_i^n(x), u_i^n(x)\rangle$ where $u_i^n(x)$ spans the invariant one-dimensional bundle given by $[26, Prop.\,3.2]$ for $\frac{dx}{dt} = A^n(x)x$. Since $A_i^n(x) \to A_i(x)$ uniformly for $x \in N_1((A))$ as $n \to \infty$, $[26, Prop.\,3.2]$ implies that $u_i^n(x) \to u_i(x)$ converges uniformly for $x \in N_1(\Lambda(A))$. Given $\in > 0$, choose $n$ sufficiently large so that $|h_i^n(x) - h_i(x)| \leq \in$ for all $x \in N_1(\Lambda(A))$ and $|\int_c h_i^n(x)\,d\mu\,(x) - \int_c h_i^n(x)\,d\mu_n\,(x)| \leq \in$.

Then

$\lambda_i(\mu) \leq \int_c |h_i(x) - h_i^n(x)|d\mu\,(x) + |\int_c h_i^n(x)\,d\mu\,(x) - \int_c h_i^n(x)\,d\mu_n\,(x)|$

$+ \int_c h_i^n(x)\,d\mu_n(x)$

$\leq \quad 2\in$

Since $\in > 0$ is arbitrary, $\lambda_i(\mu) \leq 0$ for all $i$ and $\mu$ is saturated.

Using this lemma, we prove that if K is saturated, then there exists a saturated invariant measure $\mu$ supported by K. Assume K is saturated. Work of Hofbaue and So $[7, Thm.\,2.1]$ implies that either there exists $y \in W^s(K) \cap C^+$ or K is not isolated for the unrestricted flow $x.t$ If there exists $y \in W^s(K) \cap C^+$, then for all $t \leq 0$ define $v_t = \frac{1}{t}\int_0^t \delta_{y.s}\,ds$ where $\delta_{y.s}$ denotes a Dirac measure based at the point $y.s$. Dissipativeness of (2) and weak* compactness of the Borel probability measures supported by (A) imply there exist $t_k \to \infty$ such that the sequence $v_{tk}$ converges in the weak* topology to a Borel probability measure $\mu$ with support in K.A standard argument

implies that $\mu$ is $x.t$ invariant. Define $h_i: C \to R$ by $h_i(x) = (A_i(x)u_i(x), u_i(x))$.Proposition 2 and weak* convergence imply that

$$0 \geq (\lambda i(y) = \limsup_{t\to\infty} \frac{I}{t}\int_0^t h_i(y.s)ds$$

$$= \limsup_{t\to\infty} \int_c h_i \, dv_t$$

$$\geq \limsup_{t\to\infty} \int_p h_i \, d\, v_{tk} = \lambda i(\mu)$$

For all $i$. Alternatively suppose that K is not isolated for the semi flow Then there exists a sequence of positive $\omega - limit$ sets that accumulate on K. Let $\mu_n$ be a sequence of ergodic probability measures supported by these $\omega -$ limit sets. Propositions 1 and 3 imply that $\mu_n$ are saturated for all $n$ Applying Lemma I with $\delta_n$=0 for all $n$ implies that there exists a saturated invariant measure $\mu$ suppoorted by *K*. Hence we have shown that if *K* is saturated then there exists a saturated invariant probability measure $\mu$ supported by *K* Equivalently, (4) holding for all invariant probability measures with support in *K* implies that *K* is unsaturated.

Next we show that if *K* not robustly saturated, then there exists a saturated invariant measure supported by *K* Indeed suppose *K* is not robustly saturated, Then there exists a non-negative sequence $\delta_n$ converging to zero and a sequence of saturated measures $\mu_n$ for $\delta_n$-perturbations of (2)with support in the continuation of *K*. Let $\mu$ be a weak* limit point of $\{\mu_n\}\infty_{n=1}$ Lemma 1 implies that $\mu$ is saturated. Moreover, since the continuation of *K* converges to *K* as $\delta_n \to 0, \mu$ is a saturated invariant measure for (2) supported by *K* Hence we have shown that if *K* is not robustly saturated then there exists a saturated invariant measure supported by *K* Equivalently if (4) holds for all invariant measures with support in *K* then *K* is robustly unsaturated.

To see the equivalence of the conditions given by (4) and (5) let $\Delta = \{p \in R_+^k : \sum_i P_i = 1\}$ and notice that

$$\text{Min max } \lambda i \, (\mu = \min_{\mu} \max_{p\epsilon\Delta} \sum_i Pi\lambda i(\mu)$$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K* The Minimax theorem (see, e. g.,[27] implies that

(6) $\min \max \lambda i(\mu) = \min_{\mu} \max_{p \epsilon \Delta} \sum_i Pi\lambda i(\mu)$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K*. Since min$\mu$ $\sum_i P_i \lambda_i(\mu)$ is attained at an ergodic probability measure with support in *K*, the equivalence of the conditions given by (4) and (5) is established.

To prove the final assertion of the Theorem assume $x \to A(x)$ is twice continuously differentiable and there exists a saturated ergodic measure $\mu$ supported by *K* Proposition 3 implies that there exists *I*C {1,…,k} such that (□$i \in iC_i^+$)=1and $\lambda i(\mu) = 0$ For all $i \in I. SinceKc\,.\partial$□,$\{1, \ldots . k\}\backslash I$ is non-empty. We will show that (2) is not robustly permanent by proving that for all $\delta > 0$ there exists a $\delta -$perturbation of (2) for which the continuation of *K* is saturated Let $\delta > 0$ be given Choose V $c$ N$\delta$ ($\Gamma$(A)) to be a compact neighborhood of Γ(A) such that V.t $c$ int *V* for all t > 0.Let $\eta > 0$ be such that *W*=:*N* $\eta$(Γ(A))$c$ *V* Let p:c→ $[0,1] be$ a$C^{\infty}$function such that $p(x) = 1 \quad for\ all\ x \in \Gamma(A) and\ p(x) = 0 for\ all\ x \in C\backslash W\ Define\ \tilde{A} = (\tilde{A}_1, \ldots, \tilde{A}_k)$ by

$$\tilde{A}_i(x) = \begin{cases} \frac{A(x)}{A(x) - \frac{\delta}{2Idp\,(x)}} & if\ i \epsilon I \\ if\ i \notin I \end{cases}$$

Where Id denotes the identity matrix of appropriate dimension Let x. t denote the semi flow of $x = \tilde{A}(x)x$.Since$x.t = x.t$ $whenever\ x.[0,t] \in C\backslash W, it\ follows\ that\ \ \tilde{V}.t \subset intV\ for\ all\ t > 0. Hence, \Gamma(\tilde{A})c\ V\ c\ N\delta\ (\ (A))). We\ \ also\ have\ ||A(x) - \tilde{A}(x)|| \leq \frac{\delta}{2}$Therefore $x = \tilde{A}(x) is\ a\ \delta -$ perturbation of (2) By construcation,$x.t = x.t\ for\ all\ x \in K and\ t \geq 0. Hence, \mu$ is ergodic for the semiflow of $\dot{x} = \tilde{A}(x)x\ and\ \lambda_i(\mu) \leq -\frac{\delta}{2}$ for this semiflow and $i \in \{1, \ldots, k\}\backslash I$.Let $\mathcal{L}$ and O be the Lyapunov exponents and Oseledec regular points supported by $\mu$ for $x = \tilde{A}(x)x(see, e.g.,$

[28] for definitions).At each point $x \in O$, the splitting of $R^n$ determines three subspaces: the stable subspace $E^s(x)$, the center subspace $E^c(x)$, and the unstable subspace $E^u(x)$ Proposition 3 and our choice of Ã imply that $E^s(x) \cap C^+ \neq \emptyset$. The Pesin stable manifold theorem [28, Corollaries 3.17 and 3.18] implies that tangent to $E^s(x)$ $E^c(x)$ and $E^u(x)$ are locally $x.t-$invariant families of $C^1$ discs $W^s_x, W^c_x$ and $W^u_x$ corresponding to the stable ,center and unstable manifolds The family of stable manifolds $W^s_x$ is contained in $W^s$ (Ã,supp $(\mu)$). $Since\ E^s(x) \cap C^+ \neq \emptyset$, $W^s_x \cap C^+ \neq \emptyset$ for some $x \in K$. Consequently $K$ is saturated for $x = \tilde{A}(x)x$.

The difference between the sufficient and necessary condition is that for the sufficient condition (4) has to hold for all invariant measures supported by *K*, while for the necessary condition (4) has to hold only for ergodic measures supported by *K*. since the invariant measures lie in the convex hull of the ergodic measures the sufficient condition can be more restrictive then the necessary condition.

For three Classes of invariant sets the necessary and sufficient conditions coincide Recall a compact invariant set *K* for (2) is uniquely ergodic if K only supports one invariant measure Recall a compact invariant set *K* is Axiom A if the flow of (2) restricted to *K* is transitive and hyperbolic (see, e.g.,[29] for definitions.)

**Corollary 1:** Assume $x \to A(x)$ is twice continuously differentiable It K$\subset \partial$□ is a compact invariant set and either is uniquely ergodic or Axiom A, then *K* is robustly unsaturated if and only if (4) holds for all ergodic, measures $\mu$ with support in *K*

Proof. If K is uniquely ergodic then the assertion follows immediately For an Axiom A invariant set Sigmund [30, thm.1] has proven that invariant measures supported by periodic orbits of *K* are dense in the set of invariant probability measures supported by *K* In particular any invariant measure can be approximated by an ergodic measure and consequently the result follows.

Another special case where the necessary and sufficient conditions coincide is when *K* supports all populations except one for discrete time models this case was considered by [22].

**Corollary2:** Assume $x \to A(x)$is twice continuously differentiable Let K$c\ II_i \geq 2C_i^+$ be a compact invariant set. Then K is robustly unsaturated if and only if $\lambda_1(\mu) > 0$ for all invariant measures supported by *K*

## 4. MORES DECOMPOSITIONS AND ROBUST PERMANENCE

To state the sufficient condition for robust permanence, we use a characterization of permanence due to Garay [8] and Hofbauer and So [7] that involves Morse decompositions of the boundary flow Conley[31] defined a collection of sets$\{M_{1,\ldots,}M_m\}$ to be a Morse decomposition for a compact invariant set $K$ if $M_1\ldots,M_m$ are pairwise disjoint, compact isolated invariant sent for the flow of (2) restricted to K. For each $x \in K$ there are integers r=r$(x) \leq s = s(x)$ such that $a(x) \subseteq M_r$and $\omega(x) \subseteq M_{s.}$ If $r(x) = s(x), then\ x \in M_{r.}$ Garay, Hofbauer and so [7,8] proved the following characterization of permamnce

**Theorem 2** (Garay, Hofbauer-so) if $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$?then (2) is permanent if and only if each of the components $M_1$ are unsaturated.

Theorem 1 and 2 imply the following result: As stated in the above all these results are true for most of problems, interactions, transactions in cosmology, econometrics, finance and other subjects, which we shall illustrate as we go on constructing models for the same.

**Theorem 3:** If $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$?and (4) holds for each of the components of the More decomposition, then (2) is robustly permanent Conversely if $x \to A(x) is\ twice\ continuously\ differentiable$

*and (4) is violated by an ergodic measure supported by one of the components of the Morse decomposition, then (2) is not robustly permanent.* Theorem 3 in conjunction with Corollaries 1 and 2 yield a characterization of robust permanence for a class of structured models.

**Corollary 3:** *Let $\{M_1,\ldots..,M_m\}$ os a Morse decomposition for* $\partial$C. *Assume x* $\mapsto$ *A(x) is twise continuously differentiable and for each Morse component $M_i$ one of the following assertions hold $M_i$ is Axiom A, $M_i$ is uniquely ergodic, or There exists j* $\in$ { 1,….,k} *such that $M_i$ ----- { $x \in C : x_l \gg 0$ for $l \neq j$} Then (2) is robustly permanent if and only if (4) holds for all ergodic measures supported by $U_iM_i$.*

Our result also provided a structured analogue for characterizing totally permanent systems[32]. For an ergodic probability measure$\mu$, define sup ($\mu$) to be the subset $I$ -----{1,…,k} such that $\mu(---_i \in_I C_+) = 1$

**Corollary 4:** *The following statements are equivalent: For all ergodic probability measures μ with support in ∂C, $\lambda_i(\mu)>0$ for all $I \in \{1,\ldots,k\} \setminus supp(\mu)$ (2) and all of its subsystems are robustly permanent.*

*Proof. Suppose the first statement holds. Then unsaturated condition given by (5) with p= (1,…,1) and Theorem 3 implies that (2) and all its subsystems are robustly permanent.. The other direction follows form the necessary condition for robust permanence in Theorem 3 and the second assertion of Proposition 3.*

*To illustrate the broad applicability of theorem 3, we develop applications to spatially structured models in section 5 and to a disease model and a gene network model in section 6.*

## 5. APPLICATIONS: PATCH MODELS

A fundamental application of our result are to spatially structured models, with *k* species dispersing between *m* patches {33-35}:

(7) $$\frac{dx_i^j}{dt} = x_i^j f_i^j(x^j) + \sum_l d_i^{jl} x_i^l - e_i^j x_i^j \qquad i = 1,..,k \qquad j=1,\ldots,m$$

Where $x_i^j$ denotes the density of species I in patch *j*, $x^j = (x^j_1,\ldots..,x^j_k)$ is the vector of species densities in patch *j*, $f^j_i$ is the per-capita growth rate of species *i* in patch *j*, $d_i^{jl} \geq 0$ is the dispersal rate for species *i* from patch *l* into patch *j*, and $e_i^j$ is the emigration rate of species *i* out of patch *j*. Hence $e_i^j \geq \sum_l d_i^{lj}$. We assume that the matrices $(d_i^{lj})_{j,l}$ are irreducible for each species *i* . it is then easy to write (7) in the form (2), with $x_i = (x^l_i,\ldots, x_i^m)$ 'where' denotes transpose.

For a single species (7) generates a monotone flow. Under mild assumption, e.g., each $f^j_i$ *is* decreasing and negative for large densities, there is a globally stable equilibrium, see, e.g.,[36] and [37, sect. 5.4]. This equilibrium will be positive if the invasion rate at the origin λ (0), which is given by the leading eigenvalue (stability modulus) of the matrix $(d^{jl} + f^j(0) - e^j)$, is positive.

5.1 Two types of debitable and creditable systems:. For *two competing species*, e.g., $f^j_i$ is decreasing with respect to both species and is negative at large densities, the dynamics are still monotone, and hence almost all orbits converge

to an equilibrium. Robust permanence requires two conditions. First, $\lambda_i$ (0) >0 for both species in which case there are two single species equilibria $E_1$ and $E_2$. Second, the invasion rates $\lambda_2$ ($E_1$ ) and $\lambda_2$ ($E_2$) are positive. Whether there is a unique positive globally stable equilibrium for the robustly permanent system depends in a delicate way on the system parameters, see [38]. For slowly dispersing populations (i.e. $0< e^j_i \ll 1$ for all *i, j* ), the robust permanence condition is particularly

Straight forward to verify. Let $E^j_i$ be the largest solution to $x^j_i \quad f_i{}^j$ ($x^j_i$) = 0 i.e. the equilibrium attained by species *I* in patch *j* when there is no dispersal and no competitors. For slowly dispersing populations, $E_1$ is close to the product ($E_1^1$ ,…,$E_1^m$), and the matrix $A^2$ ($E_1$) is close to the diagonal matrix diag($f^i_2(E_1{}^j)$)$^m_{j=1.}$ Hence the invasion rate $\lambda_2$ ($E_1$) is close to $\max_j f_2^j (E^j_1)>0$. So we obtain: The system with two slowly dispersing competing species is robustly permanent if and only if system refers to system of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

(8) $\max_j f_2^j (E^j_1)>0$ and $\max_j f_2^j (E^j_2)>0$

Intuitively, robust permanence requires for each species there is at least one patch where it can persist.

For *systems, particle interactions, financial transactions, cosmological interactions and reactions, as we shall show model by model the* situation are even easier. Since (under mild assumption on the prey dynamics) the Morse decomposition of $\partial$C is given by two equilibria {0,$E_1$}, the spatial predator-prey system is robustly permanent, if the predator can invade at the prey equilibrium: $\lambda_2$ ($E_1$) > 0. However, the global dynamic are likely to be more complicated. For instance, if in each patch there is a globally stable limit circle (e.g.,[39-41]) and dispersal is sufficiently slow (i.e. $0< e^j{}_i \ll 1$for all *i, j*), our result about robust permanence and the theory of normally hyperbolic manifolds [42] imply that there is a positive *m* dimensional torus, which attracts almost all orbits in C, whenever $\max_j f_2^j (E^j_1)>0$

**5.2 Rock-Paper-Scissors Dynamics.** The Lotka-Volterra model of rock-paper-scissor dynamics is a simple model that is used as property for understanding intransitive ecological outcomes [43,44]. Here, a simple spatial version of this dynamic is given by

$\frac{dx_1^j}{dt} = x_1^j(1 - x_1^j - \beta^j\ x_2^j\ - \alpha^j\ x_3^j\ ) + \sum d_1^{jk}\ \ x_1^k\ - e_1^j\ x_1^j$

$\frac{dx_2^j}{dt} = x_2^j(1 - \alpha^j\ x_1^j\ - x_2^j - \beta^j\ \ x_3^j\ ) + \sum d_2^{jk}\ \ x_2^k\ - e_2^j\ x_2^j$

$\frac{dx_3^j}{dt} = x_i^j(1 - \beta^j\ x_1^j\ - \alpha^j x_2^j - \beta^j\ \ x_3^j\ ) + \sum d_3^{jk}\ \ x_3^k\ - e_3^j\ x_3^j$

Where $a^j \in (0,1)$ and $\beta^j > 1$ for all $j$. A more general version of this dynamic is presented in [44]. Under the assumption that there is no cost to dispersal (i.e. $e^j_i = \sum_l d^{l,j}_i$ for all $i, j$) , the maximal invariant set in $\partial C$ consist of the origin 0 and a heteroclinic cycle connecting positive single species equilibria$E_1$, $E_2$ and $E_3$. For slowly dispersing population ($0 < e^j_i \ll 1$ for all $i$, j), $E_1$ is close to $(x_1, x_2, x_3) = (1,\ldots,1,0,\ldots,0,0,\ldots,0)$ and the invasion rates $\lambda_2\ (E_1)$and $\lambda_3\ (E_1)$are close to $\max_j 1 - a^j$ and $\max_j 1 - \beta^j$. For the other equilibria $E_i$, the debitable rates of the missing species are also given by $\max_j 1 - a^j$ and $\max_j 1 - \beta^j$. Consider the Morse decomposition of $\partial C$ given by 0 and the heteroclinic cycle. Since $\lambda i\ (0) = 1 > 0$ for $i = 1, 2, 3, 0$ is robustly insaturated. An algebraic computation reveals that (4) holds for all invariant supported by the heteroclinic cycle (i.e. all convex combination of the Dirac measures supported by the equilibria) if and only if

$$\max_j(1 - a^j) > -\max_j\ (1 - \beta^j)$$

Equivalently,

$$2 > \min_j a^j + \min_j \beta^j$$

In particular, even if the heteroclinic cycle is attracting for each patch when the system is uncoupled (i.e. $2 < a^j + \beta^j$ for all $j$ ), it can be repelling for the weakly coupled system.

5.3. Three species **Lotka Volterra in spatially homogenous environments.** Consider Lotka-Volterra dynamics in a spatially homogenous environment: $(x + a)^n =$

(9) $$\frac{dx_i^j}{dt} = x_i^j\left(r_i - \sum_s a_{is}x_s^j\right) + \sum_l d_i^{jl}x_i^l - e_i^j x_i^j, \qquad i = 1,\dots,k;\ j = 1,\dots,m$$

We assume that dissipative can be shown by a linear Liapunov function: There are $c_i > 0$ such that $\sum_i c_i a_{is}x_i x_s < 0$ holds for all $x \in R_+^k \setminus \{0\}$. Then the weighted sum of densities across all patches $\sum_{i,j} c_i x_i^j$ is decreasing for large densitities and hence (9) generates a dissipative semiflow with global attractor $r$. We further assume that for each patch *j* and each species *i* the immigration rate equals the emigration rate , i. e,

(10) $$e_i^j = \sum_l d_i^{jl}$$

This guarantees that the set $H = \{x \in C: x_i^j = x_i^l \ \forall i,j,l\}$ of spatially homogenous states is forward invariant under (9).

Now consider the Lotka-Volterra dynamics without spatial structure

(11) $$\frac{dx_i}{dt} = x_i\left(r_i - \sum_s a_{is}x_s\right), \qquad i = 1,\dots,k$$

(11) $$\frac{dx_i}{dt} = x_i(r_i - a_{is}x_s), \quad i = 1,\dots,k$$

It has been shown [12,ch.16.1,16.2] That for *k* =3,(11) is robustly permanent (robust meaning here within the class of Lotka –Volterra systems) if and only if all equilibria on the boundary of $R_+^3$ are unsaturated and whenever there is a heteroclinic cycle connecting the one species equilibria then this cycle is repelling. Equivalently the following four conditions hold. There exists an interior equilibrium $\dot{x}$ (i.e $A\dot{x} = r \ \ with\ \dot{x} \gg 0$); Det (-A)>0; the 2 species subsystems are not bistable competition systems if there is a heteroclinic cycle between the one species equilibria, say $\mathbf{E_1 \to E_2 \to E_3 \to E_1}$ then the following inequality holds

12) $$\lambda_2(E_1)\lambda_3(E_2)\lambda_1(E_3) > |\lambda_3(E_1)\lambda_1(E_2)\lambda_2(E_3)|$$

Furthermore, the proof in [12, ch.16.1, 16.2] shows that the boundary flow of (11) has a simple Morse decomposition: if (iv) applies then there are two Morse sets, the origin 0 is a repeller, and the heteroclinic cycle is the dual attractor (within$\partial R^3_+$); if (iv) does not apply then the finitely many boundary equilibria form a Morse decomposition of the boundary flow.

**Species, population we mean types of and population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.**

**Theorem 4:** Under the above assumptions,(9) for k=3 is permanent if and only if *(i),(ii),(iii)* and *(iv)* hold.

Proof: Suppose (9) is robustly permanent. Since the nonspatial system (11) is the restriction of (9) to the invariant subspace **H**, (11) must be robustly permanent and hence*(i)-(iv)* hold

Conversely, suppose that (i)-(iv) hold. We show that the above Morse decomposition of the boundary flow of (11) (now in **H**) is also Morse decomposition for (9) restricted to the maximal invariant subset of $\partial C$ which is contained in the spatial two species subsystems of (9). For this we show that in each 2 species system the orbits of (9) converge to **H** (For the one species subsystems this is obvious.) If a two species subsystem of (11) has an internal equilibrium, say $\mathbf{E_{12}}$ then by assumption (ii),this is not a saddle but a sink, and hence $a_{11}a_{22} > a_{12}a_{21}$. This implies that the 1-2 submatrix of A is VL-stable [12,ch.15.3,],and [45] implies that $\mathbf{E_{12}}$ is the global attractor for the spatially structured system (9) restricted to the first two species.

If a two species subsystem of (11) has no internal equilibrium, then a one species equilibrium say $\mathbf{E_1}$ is the global attractor. Then Lemma 2 (see below) show the corresponding result for (9).A similar proof applies to the case where the origin **0** is the global attractor of a two species subsystem.

Finally observe that under the above assumptions (9)-(10), the invasion rates $\lambda_i$ at spatially homogenous boundary equilibria are the same for (9) and (11) Since all boundary equilibria are unsaturated for (11), they are for (9) as well. The same applies in case (iv) to the heteroclinic cycle. Note that the condition (12) on the eigenvalues is equivalent to the fact that all invariant measures supported on the heteroclinic cycle are unsaturated, see [16,Ex.4.5]

Species we mean types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

**Lemma 2:** Suppose, in a two species Lotka –Volterra system (11) (k=2) all interior orbits converge to the one species equilibrium $\mathbf{E_1}$ .Then all orbits of the spatial version (9) (k=2) converge to the spatially homogeneous one species equibrium $\mathbf{E_1} \in \mathbf{H}$.

**Proof:** If the system is of predator-prey or mutualistic type then Hasting's [45] result applies again. So let us assume that the local interaction is given by the competition system

(13) $\frac{dx_1}{dt} = r_1 x_1(1 - x_1 - ax_2)$ $\frac{dx_2}{dt} = r_2 x_2(1 - \beta x_1 - x_2)$

With 0<a <1< $beta$.In this case, Kazuo Kishimoto (letter to JH, Oct 1987) has given the following argument. There exists a family of forward invariant rectangles contracting to $\mathbf{E_1}$. This shows that all interior solutions of (9) (k=2) converge to $\mathbf{E_1}$.

Theorem 4 shows that a spatial network of identical patches is permanent, if the within patch dynamics is permanent, and three species Lotka-Volterra dynamics. It is not clear how to extend this result to more than three species. The crucial step in the three species case is that permanence precludes bistable two species subsystems. The spatial version of a bistable two species system allows plenty of stable equilibria outside **H,** see [33].However, permanent four species Lotka Volterra systems may have two species bistable subsystems [12,ch.16.4].So there may be invariant sets on the boundary outside **H** which need to be unsaturated to make (9) permanent.

## 6. MORE APPLICATIONS

6.1. **Disease, particle, product, marketing, financial, econometric dynamics with density-dependent demography**. GAO and Hethcote [46] introduced a model of disease dynamics with density-dependent demography. Here we consider a variation of their model in which a population of size P has a per-capita birth rate b(P) and per –capita death rate d. To ensure that the population persists in the absence of the disease, we assume that there exists K>0 such that b(K)=d. To describe the disease dynamics, let *S* be the number of individuals susceptible to the disease, *I* the Number of infected individuals, and R the number of individuals that have recovered from the disease. We assume that P=S+*I* +R. On the same lines we can think of sick and non performance accounts, and those which are conducting themselves well, and those which are satisfactory and find a parallel formula. The same thing

holds true in case of a black hole eating matter and black hole radiating radiation.. If the disease transmission is asymptotic, β is the contact rate between susceptible and infected individuals, γ is the rate at which individuals recover from the disease, and m the mortality rate due to disease, then the population dynamics is given by

$$\frac{dp}{dt} = b(P)P - dP - mI$$

$$\frac{dI}{dt} = \beta(P - I - R)\frac{I}{\epsilon + P} - (\gamma + d + m)I$$

$$\frac{dR}{dt} = \gamma I - dR$$

Where $\epsilon>0$ is a constant It is useful to introduce a change of coordinates in which $y=\frac{I}{P}$ and $z=\frac{R}{P}$:

$$\frac{dP}{dt} = (b(P)-d-my)P$$

$$\frac{dy}{dt} = \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(P)\ y$$

$$\frac{dz}{dt} = \gamma y - (b(P) - my)z$$

Setting $x_1$=P and $x_2$=(y,z) 'where' denotes transpose yields a structured model where $A_1(x)$=b-d-my and

$$A_2(x) = \begin{pmatrix} \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(p) & 0 \\ \gamma & my-b(p) \end{pmatrix}$$

Since {0,(*K*,0,0)}is a Morse decomposition of the boundary dynamics Theorem 3 and Remark 1 imply that this model is robustly permanent if and only if b(0)>d and $\beta\frac{K}{\epsilon+K} > \gamma + m + d$. In particular, for $\epsilon > 0$ sufficiently small, one requires that the basic reproductive number $\frac{\beta}{r+m+d}$ is greater then one.

Remarkably this same criterion determines robust permanence of model with significantly more complicated boundary dynamics. For example, suppose the focal population is a predator species. If the prey has abundance *N* and exhibits logistic dynamics f(*N*)is the per-capita predator consumption rate of the prey and b(N)is the per-capita reproductive rate of the predator then the dynamics become

$$\frac{dN}{dt} = rN(1 - N/K - f(N)P$$

$$\frac{dP}{dt} = (b(N) - d - my)P$$

$$\frac{dy}{dt} = \left[ \beta(1 - y - z)\frac{P}{\epsilon + P} - r - m(1 - y) - b(N)\, y \right]$$

$$\frac{dz}{dt} = \gamma y - (b(N) - my)z$$

Where r>0 and *K*>0, and *f*(*N*),b(N)satisfy f(0)=b(0)Define $x_1$=*N*, $x_2$=*p*, $x_3$=*(y,z)*′ Let us assume that the disease-free community (i.e.y=z=0)is permanent (i.e.b(K)>d) and let A be the global attractor in the interior of the P-N plane Under these assumption a Morse decomposition of the boundary dynamics is given by $(x_1, x_2, x_3) = (0,0,0)$, $(x_1, x_2, x_3) = (K, 0,0)$and A. Since the equilibria (0,0,0) and (k,0,0) are unsaturated, it remains to characterize $\lambda_3(\mu)$for any invariant measure μ supported by A. Proposition 3 and Remark 1 imply

$$\lambda_2(\mu) = \int_c (bx_1)\text{-}d)d\mu(x) = 0$$

*Hence* $= \int_c bx_1)d\mu(x) = d$ *and*

$$\lambda_3(\mu) = \beta \int \frac{x_2}{\epsilon + x_2} d\mu(x) - r - m - d$$

**Theorem 3**: and Remark 1 implies that this structured model is robustly permanent if $\frac{\beta}{r+m+d} > 1 and\ \epsilon > 0$ is sufficiently small.

6.2. Gene networks. The repressilator is an oscillatory gene network based on three (or more generally an odd number of transcriptional repressors. Such a system has been genetically engineered in E-coli by Elowitz and Leibler [47] Mathematical model go back to [48.49] they all involve concentrations of proteins and mRNs. A modified model with auto-activation was suggested in [50]

$$(14)\ x_i = \beta(y_i - x_i)$$

$$(15)\ y_i = aF(x_{)i} \text{-} y_i)$$

With

$$(16)\quad F(x)_i = x_i g\,(x_i, x_i - 1) = \frac{x_i}{1 + x_i + p x_{i-1} + Kp\, x_i x_i - 1}$$

Here $x_i$ and $y_i$ are normalized concentrations of proteins and mRN As belonging to gene $i$,p is the strength of repression and k a cooperativity parameter this is a structured system of type(2)

$$(17)\ \begin{pmatrix} x_i \\ y_i \end{pmatrix} dot = \begin{pmatrix} -\beta & \beta \\ ag(x_i,x_i) & -1 \end{pmatrix} \begin{pmatrix} x_i \\ y_i \end{pmatrix}$$

In [50,Thm 3] it was shown that for n=3,a>1 and p>1,system(14)has a heteroclinic cycle connecting the 3 single gene equilibria (similar to the rock –scissors-paper dynamics of section5.2)Moreover

- the system is permanent i.e. the boundary of $R^6$+is repelling, if $\lambda + \mu > 0$
- this heteroclinic cycle is asymptotically stable if $\lambda + \mu < 0$

where $\lambda > 0$ and $\mu < 0$ are the invasion rates of the two missing genes i.e. the leading eignvalues of the 2X2 matrix in (17)evaluated at the appropriate single gene equilibria (for explicit expressions see[50,eqs.(127,128)].)

Theorem 3 implies the permanence condition implies the stronger conclusion of robust permanence A similar result holds for n odd, but the heteroclinic cycle is then actually a heteroclinic network.

Note: For references please see the original paper.

One of the examples for which we apply the model are as follows: **Progress in Particle and Nuclear Physics Volume 88, May 2016, Pages 1–28 Chiral magnetic and vortical effects in high-energy nuclear collisions—A status report D.E. Kharzeeva, b, J. Liaoc, d, , , S.A. Voloshine, G. Wangf**

The interplay of quantum anomalies with magnetic field and vorticity results in a variety of novel non-dissipative transport phenomena in systems with chiral fermions, including the quark–gluon plasma. Among them is the Chiral Magnetic Effect (CME)—the generation of electric current along an external magnetic field induced by chirality imbalance. Because the chirality imbalance is related to the global topology of gauge fields, the CME current is topologically protected and hence non-dissipative even in the presence of strong interactions. As a result, the CME and related quantum phenomena affect the hydrodynamical and transport behavior of strongly coupled quark–gluon plasma, and can be studied in relativistic heavy ion collisions where strong magnetic fields are created by the colliding ions. Evidence for the CME and related phenomena has been reported by the STAR Collaboration at Relativistic Heavy Ion Collider at BNL, and by the ALICE Collaboration at the Large Hadron Collider at CERN. The goal of the present review is to provide an elementary introduction into the physics of anomalous chiral effects, to describe the current status of experimental studies in heavy ion physics, and to outline the future work, both in experiment and theory, needed to eliminate the existing uncertainties in the interpretation of the data. Keywords

Chiral magnetic effect; Chiral vortical effect; Chiral anomaly; Quark-gluon plasma; Heavy ion collisions Corresponding author at: Physics Department and Center for Exploration of Energy and Matter, Indiana University, 727 E Third Street, Bloomington, IN 47405, USA. Copyright © 2016 Elsevier B.V. All rights reserved.

## Sufficient Conditions For The Existence And Linear Stability Of Time Independent Age Distributions In Respect Of Particles, Propositions, Problems, Hegelian Structures, Deleuzean Treatises, Interactions And Transactions In Cosmos And Others

**Abstract:** Gopalaswamy studied sufficient conditions are obtained for the existence and linear stability of time independent age distributions in two species competition with age time lagged density dependent mortality and fertility functions.. Work was published in Bulletin Australian Mathematical Society Volume 25 (1982). We show how the model could be applied in case of finance, econometrics, psycho economics, interactions between particles and transactions in nature. Basically the competition is for resources, notwithstanding such experiments could be conducted in laboratory like a rat kept in front of a cat, such incidents could be represented mathematically, and the corresponding linear stability conditions ca be obtained in case of interactions in cosmos. Take for example the case of a cash credit account of a borrower. Many inward bills waiting for paid, or inland or foreign bills payment all are in line for the debit of the account, which also acts as a resource. In addition there are recurrent cases where in the credit is issued to the cash credit account in case of demand and inland bills. In case of economics, investment has inventories and it has to be utilized towards the end of consummation of output. And this is case of output gobbling part of investment which is inventories. Output has to be sold to the customer and the amount received is diverted partly to savings and consumption and part could be employed towards the consolidation of the system. So this is a case of investment itself acting as a both receiver and payer. It shall eat and also produce the finance. Savings as stated above could be deployed for the further diversification and proliferation and that is towards the end of putting the company on terra firma. Another example in cosmology is the case of black holes guzzling matter. There are lot of black holes and lot of matter. Result of such a guzzling the production of Hawking radiation. Horava gravity is another example. How the conformal field theory could be changed in the case of Horava gravity. Still another case is the case of Vonneumann algebra in the cae of Horava gravity. In this case like cash credit account, the Horava gravity and Von Neumann algebra act as both receivers and takers. Such incidents are legion in cosmology and even in philosophy. We have considered self as individual general ledger, and the debits and credits are raised to the same account thus anointing as both a receiver and payer. We give plethora of such cases and apply the model. Due to spatial constants, same module is used for all the modules and we rewrite to state that the model is applicable to all modules. This is due to spatial restrictions. What is to be done is that each and every parameter is to be changed with a subscript or a superscript and finally it is to be concatenated towards the end of obtention of prediction model or optimizing model, for Duality of motion model for further analysis. Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation

of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

**Introduction**

Let $P_1$ *(t)* and $P_2$ *(t)* denote the total particles, interactions, transactions, etc, which we dilate upon further investigation. Perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. **Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems** sizes (or biomasses) at time $t\geq 0$ of two interacting species living in a common habitat and competing for a common pool of resources. The competition for resources will be implicit in our model similar to that in the two species Lotka- Volterra competition system. Assuming constant sex ratios in the two species we can consider $P_1$ *(t)* and $P_2$ *(t)*to be the Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

population of f; immigration, emigration and internal dispersion in the habitat are assumed to play no significant role in the dynamics of the community.Popoulation we mean Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

We suppose that the two species contain respectively $p_1\,(a, t)\,da$ and $p_2\,(a, t)\,da$ individuals with ages between $a$ and $a + da\ (a\geq 0)$ at time $t$ so that we have

(1.1) $P_1\,(t) = \int_0^\infty \rho_1(a, t)\,da,\ P_2\,(t) = \int_0^\infty \rho_2(a, t)\,da,\ t\geq 0.$

The rates of change of the two species densities are by definition $D\rho_1\,(a, t)$ and $D\rho_2\,(a, t)$ where

(1.2) $D\rho_i\,(a, t) = \limsup_{h\to 0} \frac{\rho_i(a+h,t+h)-\rho_i(a,t)}{h},\ i = 1, 2.$

Assuming the existence of the limits in (1.2) we consider the following time lagged model system

(1.3) $$\begin{cases} D\rho_1\,(a,t) + f_1\,(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_1(a, t) = 0, & a>0,\ t>0 \\ D\rho_2(a,t) + f_2(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_2(a, t) = 0,' \end{cases}$$ (1.4) $$\begin{cases} \rho_1(0, t) = \int_0^\infty b_1\,(a, P_1\,(t), P_2\,(t-\tau))\rho_2(a, t)da\,, & t>0, \\ \rho_2(0, t) = \int_0^\infty b_2\,(a, P_1\,(t), P_2\,(t-\tau))\rho_1(a, t)da\,, \end{cases}$$ $\rho$

(1.5) $\rho_i\,(a, s) = \varphi_i\,(a, s),\ a.0,\ s\epsilon\,[-\tau, 0],\qquad i = 1, 2,$

(1.6) $P_i(s) = \int_0^\infty \varphi_i\,(a, s)\,da,\ s\in\,[-\tau, 0],$

Where $\tau$some is fixed nonnegative constant; $f_1$ and $f_2$ denote the age and density dependent mortality rates with time lags in interspecific interactions, $b_1$ and $b_2$ denote the age and density dependent fertility functions again having time lags in the interspecific interactions; $\varphi_1$ and $\varphi_2$ denote the initial age distributions needed for the formulation of the model. Note we mean the population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

The model system (1.1)-(1.6) will be meaningful only if $\varphi_1$, $\varphi_2$, $f_1$, $f_2$, $b_1$, $b2$ are nonnegative functions of the respective arguments. Also since the total initial populations have o be finite, $\varphi_1(\cdot, s)$ should belong to $L_1(R^+)$ for each $s\epsilon$ $[-\tau, 0]$. That $f_i$ and $b_i$ $(i = 1, 2)$ are independent of t explicitly and this assumption can be interpreted to represent the temporal uniformity of the environment.

One of the fundamental questions for (1.1)-(1.6) is the following; under what conditions on $f_i$, $b_i$ and $\tau$ does the system (1.1)-(1.4) have time independent nontrivial solutions, and when such solutions, and when such solutions exist, are they stable with respect to some suitable stability criterion? When such stable time invariant solutions exist for (1.1)-(1.4) we say that the two types has a stable age distribution. Existence of a stable age distribution is the analogue of the existence of a stable steady state for the corresponding age independent system. Types we mean types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

In this article we investigate the above question for the case of a competitive interaction with some additional constraints on the vital rates $f_i$ and $b_i$ $(i = 1, 2)$. Age dependent population systems without time delays have been considered by several authors (Gurt in and MacCamy [3, 4, 5], Gurt in and Levine [6], Haimovici [7] and Rotenberg [8]). The question of the relation between an age independent system and a corresponding age dependent system has been considered by Gurt in and MacCamy [4] and Gurt in and Levine [6] who have established an asymptotic (*as* $t\rightarrow\infty$ ) relation between such models by constructing a higher dimensional 1 umped parameter system (in terms of ordinary differential equations) to represent the age dependent distributed parameter system. Assuming the existence of stationary age distributions Haimovici [7] considers their stability in a system of two interacting populations explicitly taking into consideration the dynamical nature of the encironments resources and pollution.

In this article we add another realistic feature namely time lags in the age dependent models and obtain sufficient conditions for the existence of stationary age distributions in (1.1)-(1.4) and show that the time lags in interspecific

interactions have no effect on the linear stability of the age distributions although the decay rates of perturbations will depend on the magnitudes of the time delays. Such "harmless" natures of time lags in interspecific interactions have been noted by the author for the age independent models in [1, 2].

## 1. Existence and uniqueness of solutions

We first narrate our assumptions on the vital rates of $f_i$ and $b_i$ $(i = 1, 2)$ so that the system (1.1)-(1.6) will denote a competition system. Let $BC\ (0, \infty)$ denote the linear space of bounded continuous real functions on $(0, \infty)$ with the norm defined by

$$||f|| = \{max\ |\ f(t)\ |,\ t\ \in (0, \infty)\}.$$

Let $K_1$ be the cone of nonnegative functions of $BC\ [0, \infty]$. Let $C\ [-\tau, 0]$ denote the space of real continuous functions with a norm defined by $||g|| = \{\ max\ |g(s)\ |,\ s \in [-\tau, 0]\}$; let $K_2$ be the nonnegative cone of functions in $C\ [-\tau, 0]$ . We present our first set of assumptions on $f_i$, $b_i$ and $\varphi_i$ for $i = 1, 2$.

$A_1$. $f_1 \in C\ (R^+ \times K_1 \times K_2, R^+)$, $f_1 \geq d_1^*$ on $R^+ \times K_1 \times K_2$

$f_2 \in C\ (R^+ \times K_2 \times K_1, R^+)$, $f_2 \geq d_2^*$ on $R^+ \times K_2 \times K_1$ ,

( $d_1^*$ , $d_2^*$ are positive constants);

$f_1$ and $f_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ ; that is, *for* $i = 1, 2$,

$$|f_i\ (a, P_1^{'}\ , P_2^{'}) - f_i(a, P_1^{"}\ , P_2^{"})| \leq d_i\ [||P_1^{'} - P_1^{"}, P_2^{'} - P_2^{"}||]$$

($d_1$, $d_2$ are positive constants),

$$\partial f_1/\ \partial P_1, \partial f_1/\ \partial P_2 \in C(R^+ \times K_1 \times K_2, R^+)\ ,$$

$$\partial f_2/\ \partial P_1, \partial f_2/\ \partial P_2 \in C(R^+ \times K_2 \times K_1, R^+),$$

$A_2$. $b_1 \in C\ (R^+ K_1 \times\ K_2, R^+)$, $b_1 \leq \beta_1^* on\ R^+ \times K_1 \times K_2$ ,

$b_2 \in C\ (R^+ K_2 \times\ K_1, R^+)$, $b_2 \leq \beta_2^* on\ R^+ \times K_2 \times K_1$ ,

*($\beta_1^*$ , $\beta_2^*$ are positive constants);*

$b_1$ , $b_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ *; that is for* $i = 1, 2$,

$$|b_i\ (a, P_1^{'}, P_2^{'}) - b_i\ (a, P_1^{"}, P_2^{"})| \leq \beta_i\ [||P_1^{'} - P_1^{"}|| + ||P_2^{'} - P_2^{"}]$$

$B_1$ , $B_2$ Are positive constants),

$b_i(\circ,(P_1,P_2)\epsilon L^{\infty}(R^+,R^+)$ For fixed $P_1,P_2,i=1,2,$

$$\frac{\partial b_1}{\partial P_2},\frac{\partial b_1}{\partial P_2}\epsilon\ C\ R^{+}\text{x}\ k_{\ 1}\text{x}\ k_{\ 2},\text{R}^{-}),$$

$$\frac{\partial b_2}{\partial P_1},\frac{\partial b_2}{\partial P_2}\epsilon\ C\ R^{+}\text{x}\ k_{\ 2}\text{x}\ k_{\ 1},\text{R}^{-}),$$

$A_3.\ \varphi_i\ (\cdot\ ,\text{s})\in \text{L}^{\ 2}\ (\text{R}^{+},\text{R}^{+})\ \cap \text{C}(\text{R}^{+},\text{R}^{+})$ for each $s\in\ [-\tau,0],$

$i,\ =1,\ 2,$

By a solution of $(1.1)-(1.6)$ on $[0,\text{T}]$, we mean

$\text{P}_i\text{: R}^{+}\ \square[-\tau,\text{T}]\rightarrow \text{R}^{+}\ \ (i,\ =1,\ 2,)$ With the following properties:

$S_1.$ $\text{P}_i(\cdot\ ,t)\in \text{L}^{\ 1}\ (\text{R}^{+},\text{R}^{+})$ for each $t\ \in[-\tau,\text{T}];$

$S_2.$ $\text{P}_i(a,\cdot\ )\in \text{C}\ ([-\tau,\text{T}],\text{R}^{+})$ for each $a\ \in \text{R}^{+}$;

$S_3.$ $\text{DP}_i$ Exists along the characteristics $t-a=$ constant

On $\text{R}^{+}\ \square\{\text{R}^{+}\cap\ [0,\text{T}]\}$ And is continuous $\ (i,\ =1,\ 2,);$

$S_4.$ $\text{P}_i$ Satisfies the system (1.1)- (1.6) for

$(a,t,)\ \in \text{R}^{+}\ \square[-\tau,\text{T}]$ .

Let us first convert the system (1.1)- (1.6) into an equivalent system of integral equations. For brevity we denote by P the pair ( $\text{P}_1,\text{P}_2$) If we let

$$\text{P}_i(a+x,t+x)=\ \ \text{P}_i(x),(i,\ =1,\ 2,)\ ,$$

Then (1.3) considered at $(1.3)$ considered at $(a+x,\ t+x)$ becomes

$$\frac{dp_{\ i}(\text{x})}{dx}+f_i(a+x,\ P(t+x))\ \ P_i(x)=0$$

Which has a solution of the form?

$$P_i(x)=P_i(0)\ \exp\left[-\int_0^{\infty} f_{\ i}\big(a+s,P(t+S)\big)ds\right]$$

And hence

$$(2.1)\ p_{\ 1}\ (a+x,t+x)=p_{\ 1}\ (a,t)\ exp\big[-\int_o^x\ f_{\ 1}\ \big(a+s,p(t+s)\big)ds\big]$$

For $a \geq t$ *we choose (a,t)=(a-t,0)* in (1.3) and *x=t* in (2.1) so that

(2.2) $p_1(a,t) = \varphi_1(a-t,0)\exp\left[-\int_o^t f_1\left(a-t+s, p_1(s)\right)ds\right]$

$$O \leq t \leq a, i = 1,2,$$

With

$$P_1(s)=(p_1(s),p_2(s-t)),$$

$$P_2(s)=(p_1(s-t),p_2(s)),$$

For $0 \leq a < t$ we choose *(a,t)=(0,t-a)in (1.3)and x=a in (2.1)* so that

(2.3) $p_i(a,t)= B_i(t-a)exp[-\int_o^a f_1(s, p_i(t-a+s))ds]$

$$t > a \geq, i = 1,2,$$

Where

(2.4) $B_i(t) = p_i(0, t), i = 1,2, t > 0.$

Define $M_i and\ L_i\ (i = 1,2)$ as follows:

$$M_i(a,t, p_i(t)) = exp\left[-\int_o^{t-a} f_1(s, p_i(a-a+s))ds\right]$$

(2.5) $L_i\left(a,t,\ p_i(t)\right) = \exp\left[-\int_o^t f_1(\text{a+s},\ p_i(s))ds\right]$

It will now follow from (1.1),(1.3),(2.2)-2.5)that

(2.6) $p_i(t) = \int_o^t B_1(a)\, M_1(a, t, p_i(t))da + \int_o^\infty \varphi_i(a, 0)\ L_i(a,t,p_i(t)da,$

$$t > 0, i = 1,2,$$

(2.7) $p_i(s) = \Phi_i(s) = \int_o^\infty \varphi_i(a, s)da, s \in [-\tau, 0], i = 1,2,$

(2.8) $\beta_i(t) = \int_0^t \beta_i(a) b_i(t-a, P_i(t)) M_i.\left(a,t,P_i(t)\right)da$

$$+\int_0^\infty b_i\left(a+t, P_i(t)\right)\varphi_i(a,0)L_i\left(a,t,P_i(t)\right)da$$

$$i = 1, 2, t > 0$$

the equivalence of (2.6)-(2.8) with (1.1)-(1.6) is established by the following whose proof is identical to a similar result of Gurtin and MacCamy [4]; hence we omit the details of the proof of the following.

THEOREM 1 : *Suppose the system* (1.1)-(1.6) *satisfies the conditions* $A_1, A_2, A_3$ . *If* $(P_1, P_2)$ *is a solution of*

*(1.1)-(1.6) then the total population* $P_1, P_2$ *and the birth rates* $B_1, B_2$ *satisfy the integral equation (2.6)-(2.8).*

*conversely if* $P_1, P_2, B_1, B_2$ *are nonnegative continuous solution of (2.6)-(2.8) on [ 0, T] and if* $P_1, P_2$ *are*

*defined by (2.2)-(2.4) on* $R^+$ x [0, T] *then such* $P_1, P_2$ *provide a solution of (1.1)-(1.6) on* $R^+$ x [0, T] .

The following a *priori* estimates are useful to prove our existence theorem below. It will immediately follow from the bounds for $f_i$ and $b_i$ that

$$M_i(a,t,P) \leq \exp[-d_i^*(t-a)] \leq 1, if \quad t \geq a,$$

*(2.9)*

$$L_i(a,t,P) \leq \exp[-d_i^* t]$$

From (2.7)-(2.8) we derive that

$$\beta_i(t) \leq \int_o^t \beta_i^* \beta_i(a) \exp[-d_i^*(t-a)]\, da + \beta_i^* \Phi_i(0) \exp[-d_i^* t], \quad i = 1,2,$$

And hence by Gronwall's inequality,

(2.10) $$\beta_i(t) \leq \beta_i^* \Phi_i(0) \exp[\delta_i t], \ \delta_i = \beta_i^* - d_i^*$$

(2.5), (2.6) and (2.10) lead to

(2.11) $$|P_i(t) - 0_i(0)| \leq \beta(0) exp[\delta_i a - d_i^*(t-a)]\, da$$

$$+\int_0^\infty \varphi_i(a,0)|1 - exp(-d_i^* t)|da$$

$$\leq \beta_i^* \Phi_i(0) exp(\beta_i^* t)\left\{\frac{1-exp[-(d_i^*+\beta_i^*)t]}{d_i^*+\beta_i^*}\right\} + \phi(0) d_i^* t$$

$$\leq [\beta_i^* \phi(0) exp(\beta_i^* t) \phi(0) d_i^*] t \quad for \ t > 0.$$

We can now prove the following result on the existence and uniqueness of solution of the integral equations (2.)-(2.8)

**THEOREM 2:** *Assume that the system* (1.1)-(1.6) *satisfies the hypotheses*

| |
|---|
| $A_1, A_2, A_3$. *Then there exists a unique set of continuous function* $(P_1, P_2, B_1, B_2)$ such *that* |
| $P_i : R^+ \to R^+, \qquad B_i : R^+ \to R^+ \ (i = 1,2\ ),$ |
| *Which satisfy the integral equations* (2.6)-(2.8) |
| Proof. We will first show the existence of local solution on same (possibly small) interval [0, T] $\subset R^+$ such that |
| $P_i \in \mathrm{C}([0,T], R^+), \quad B_i \in \mathrm{C}([0,T], R^+)\ , i = 1, 2,$ |
| Satisfying (2.6)-(2.8) and then show the $P_i$ and $B_i$ can be continued as solution for any finite interval in $R^+$ for some fixed |
| Positive number $T$ we let |
| $C^+[-\tau, T] = \{f \in \mathrm{C}([-\tau, T], R)\,|f \geq 0\}$ |
| Let $\|.\|_T$ denote the supremum norm in $\mathrm{C}^+[-\tau, T]$. Then for any |
| $(x, y) \in \mathrm{C}^+[-\tau, T] \times \mathrm{C}^+[-\tau, T]$ we define |
| $\|(x,y)\| = \|x\|_T + \|y\|_T,$ |
| $\mathrm{C}_T = \mathrm{C}^+[-\tau, T]\ , \quad \mathrm{C}_T^2 = \mathrm{C}^+[-\tau, T] \times \mathrm{C}^+[-\tau, T]\ .$ |
| |
| Members of $\mathrm{C}_T$ will be denoted by $Z_t$ with the meaning that |
| $Z\,(t+s) = z_t(s)$ *for* $s \in [-\tau, T]$ *and* $t \in [0, T]$. Now for any fixed |
| $P_t = (P_{1t}, P_{2t}) \in \mathrm{C}_T^2$, (2.8) is a linear system of uncoupled Volterra integral equations in B = $(B_1, B_2)$ and hence (2.8) will |
| have a unique continuous solution under the assumption $A_1 - A_3$. let us denote such a solution by |
| B (t) = $B^T(P_t), \quad t \in [0, T]$ |
| Or |
| $B_1(t) = B_1^T(P_{1t}, P_{2t}), \quad B_2(t) = B_1^T(P_{1t}, P_{2t}), \ t \in [0, T]$ |
| If we supply this solution $(B_1, B_2)$ in (2.6) we find that (2.) will be satisfied if and only if $(P_{1t}, P_{2t})$ is a fixed point of the |
| operator |

$$\Pi(P_t) = \{\Pi_1(P_{1t}, P_{2t}),\ \Pi_2(P_{1t}, P_{2t})\}$$

Where $\pi_1$ *and* $\pi_2$ respectively denote the right sides of (2.6) for *i =1,2* .

It is not difficult to see from the nature of $M_i$ , $L_i$ and the hypotheses on $f_i, b_i, \varphi_i$ that

$$B: C_T^2 \to C_T^2 \text{ and } \Pi C_T^2 \to C_T^2$$

(which guarantees the nonnegativity of $P_i$ and $B_i$). We will first show that

There exist positive constants $r^*, t^*$ such that $\pi$ is a mapping of a sphere

$$S_{t*}(\Phi, r^*) = \{ (P_{1t}, P_{2t}) \in C_{t*}^2 \mid \|P_t - \Phi_t\|_{t*} \le r^* \}$$

Into itself and also is a contraction where $\Phi_t$ is defined on [-$\tau$, $t^*$]

By the following:

$$\Phi_t = (\Phi_{1t}, \Phi_{2t});$$

(2.12)

$$\Phi_{it} = \begin{cases} \int_0^\infty \varphi_i(a, s)\, da \text{ for } s \in [-\tau, 0]. \\ \int_0^\infty \varphi_i(a, 0)\, da \text{ for } s \in [0, t^*]. \end{cases}$$

We note that similar to (2.11) one gets the estimates

$$|\Pi_i(P)(t) - \Phi_i(0)| \le \{B_i^* \Phi_i(0) \exp[B_i^* t] + \Phi_i(0) d_i^*\} t$$

Which together with (2.7) and (2.12) lead to

(2.13) $$\|\Pi(P_t) - \Phi_t\|_{t*} \le \left( \sum_{j=1}^2 \|\Phi_{jt}\|_{t*} [B_j^* \exp(B_j^* t) + d_i^*] \quad t^* \right)$$

Thus if we choose,

$$r^* = T \sum_{j=1}^2 \|\Phi_{jt}\|_{t*} [B_j^* \exp(B_j^* T) + d_i^*]$$

Then $\Pi$ is a mapping of $S_{t*}(\Phi, r^*)$ into itself for all $t^* \le T$. we then will have to show that for small $t^* > 0$ , $\Pi$ is a contraction.

Let $P_t^{(1)}$ *and* $P_t^{(2)}$ be arbitrarily chosen in $S_{t*}(\Phi, r^*)$ and consider $|\Pi(P_t^{(1)}) - \Pi(P_t^{(2)})|$ for $t \le t^*$.

Corresponding to the chosen $P_t^{(1)}$ , $P_t^{(2)}$ we let $B^{(1)}(a) = B(P_T^{(1)})$ and $B^{(2)}(a) = B(P_T^{(2)})$ . Then from

(2.6),

$$| \text{[illegible]}(P^{(1)}(t)) - \text{[illegible]}(P^{(2)}(t)| \leq \left| \int_0^t B_1^{(1)}(a)\left\{M_1\left(a,t,P_t^{(1)}\right) - M_1\left(a,t,P_t^{(2)}\right)\right\}da \right|$$

$$+ \left| \int_0^t [B_1^{(1)}(a) - B_2^{(2)}(a)]\, M_1(a,t,P_t^{(2)})da \right|$$

$$+ \left| \int_0^t \varphi_1(a,0)\left[L_1\left(a,t,P_t^{(1)}\right) - L_1\left(a,t,P_t^{(2)}\right)\right]da \right|$$

$$= J_1 + J_2 + J_3 \text{ (say)}$$

using the elementary inequality

$$|e^x - e^y| \leq |x-y| \ for \ x,y \leq 0$$

We derive from (2.5) that

$$\left| M_1(a,t,\ p_t^{(1)}) - M_1(a,t,\ p_t^{(2)}) \right|$$

$$\leq \int_o^{t-a} \left| f_1(s,\ p_t^{(1)}(a+s)) - f_1(s,\ p_t^{(2)}(a+s)) \right| ds$$

$$\leq d_1 \| p_t^{(1)} - p_t^{(2)} \|_{t^*} \text{(using } A_1\text{)}$$

*Now,using (2.10),*

$$\left| J_1 \right| \leq \int_o^a \beta_1^* \Phi_1(0)\, e^{\delta 1^a} \left( \int_o^{t-a} d_1 \| P_t^{(1)} - P_t^{(2)} \|_t ds \right) da$$

$$\leq d_1 \beta_1^* \Phi_1(0)\, t^2 \| P_t^{(1)} - P_t^{(2)} \|_t e^{\delta 1^t/2}$$

In a similar way

$$\left| J_3 \right| \leq \int_o^\infty \varphi(a,o) \left[ \int_o^t \left| f_1(a+s, P_t^{(1)}(s) - f_1(a+s, P_t^{(2)}(s)) \right| ds \right] da$$

$$\leq d_1 \Phi_1(0) \| P_t^{(1)} - P_t^{(2)} \|_t$$

Now we have to estimate $\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right|$ in order to estimate $J_2$, from (2.8),

$$B_1^{(1)}(a) - B_1^{(2)}(a)$$

$$= \int_o^a \left[B_1^{(1)}(\sigma) - B_1^{(2)}(\sigma) b_1\left(a-\sigma, P_t^{(1)}(a)\right) M_1(\sigma, a, P_t^{(1)}(a) d\sigma + \int_o^a \left[B_1^{(2)}(\sigma)\right.\right.$$

$$\mathrm{X}\left[b_1(a-\sigma, P_t^{(1)}(a) M_1(\sigma, a, P_t^{(1)}(a) - b_1\left(a-\sigma, P_t^{(1)}(a) M_1\left(\sigma, a, P_t^{(2)}(a)\right)\right]\right] d\sigma$$

$$+\int_o^\infty \varphi(\sigma, o)\left[b_1\left(\sigma + a, P_t^{(1)}(a)\right) - b_1\left(\sigma + a P_t^{(2)}(a)\right)\right], L_1\left(\sigma, a, P_t^{(1)}(a)\right) d\sigma$$

$$+\int_o^\infty \varphi(\sigma, o) b_1\left(\sigma + a, P_t^{(2)}(a)\right)\left[L_1\left(\sigma, a, P_t^{(1)}(a)\right) - L_1\left(\sigma, a, P_t^{(2)}(a)\right)\right] d\sigma$$

$$= R_1 + R_2 + R_3 + R_4 \text{ (say)}.$$

By the uniform Lipschitz continuity of $b_1$ in $A_2$ we get

$$\left| b_1\left(a+\sigma, P_t^{(1)}(a)\right) - b_1\left(\sigma + a, P_t^{(2)}(a)\right) \right| \leq B_1 \, /\!/ \, P_t^{(1)} - P_t^{(2)} /\!/ \text{ for } a \, \varepsilon [o, t^*]$$

And hence we have

$$|R_3(a) + R_4(a)| \leq \beta_1 \Phi_1(0) || P_t^{(1)} - P_t^{(2)} ||_{t^*} + \Phi_1(0) \beta_1 d_1 || P_t^{(1)} - P_t^{(2)} ||_{t^*}^{*a},$$

$$|R_2(a)| = [\int_0^a [B_1^{(2)}(\sigma)\{b_1(a-\sigma, P_t^{(1)}(a)) - b_1(a-\sigma, P_t^{(2)}(a))\} M_1(\sigma, a, P_t^{(1)}(a))$$

$$+ b_1(a, -\sigma, P_t^{(2)}(a))\{M_1(\sigma, a, P_t^{(1)}(a)) - M_1(a-\sigma, P_t^{(2)}(a))\}] d\sigma]$$

$$+ b_1\left(a - 0, P_t^{(2)}(a)\right)\{M_1(\sigma, a, P_t^{(1)}(a)) - M_1(a-\sigma, P_t^{(2)}(a))\}] d\sigma]$$

$$\leq \int_0^a [\beta_1^* \Phi_1(0) \, \mathrm{e}^{\delta 1^a} \beta_1 \, /\!/ \, P_t^{(1)} - P_t^{(2)} /\!/_{t^*} + \beta_1 \mathrm{d}_1 \, /\!/ \, P_t^{(1)} - P_t^{(2)} /\!/_{t^*} (a-o)] do$$

$$\leq \beta_1^* \Phi_1(0) \, \mathrm{e}^{\delta 1^a} /\!/ \, P_t^{(1)} - P_t^{(2)} /\!/_{t^*} \{\beta_1 a + d_1 a^2 / 2\}$$

Thus

$$\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right|$$

$$\leq \int_o^a \left| B_1^{(1)}(\sigma)\text{-}B_1^{(2)}(\sigma) \right| B_1^{*d}\sigma$$

$$+ \Phi_1(0) /\!/ P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} \left\{ B_1 + B_2\, d_1\, a + B_1^{*} \left( B_2 a + d_1 a^2/2 \right) e^{\delta 1^a} \right\}$$

$$\leq \beta_1^* \int_o^a \left| B_1^{(1)}(\sigma)\text{-}B_1^{(2)}(\sigma) \right| \text{do} + \left( r^* + /\!/ \Phi_t /\!/_{t*} \right) \| P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} g_1^{(a)}$$

Where $g_1(a) \rightarrow \beta_1$ as $a \rightarrow 0$. By Gronwall's inequality it will follow that

$$\left| B_1^{(1)}(a)\text{-}B_1^{(2)}(a) \right| \leq \left( r^* + /\!/ \Phi_t /\!/_{t*} \right) /\!/ P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} \{ g_1(a) + 0(a) \}$$

Using this estimate for $R_1(a)$ we get

$$|R_2(a)| \leq \left| \int_0^a [B_1^{(2)}(\sigma)\, B_1^{(2)}(\sigma),]\, b_1(a)) - b_1(a\text{-}\sigma, P_t^{(1)}(a))\, M_1(\sigma, a, P_t^{(1)}(a)) do \right|$$

$$\left| B_1^{*} \left( r^* + /\!/ \Phi_t /\!/_{t*} \right) /\!/ P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} 0(a) as\ a \rightarrow 0 \right.$$

*Thus*

$$\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right| \leq (r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ P_t^{(1)} - P_t^{(2)} /\!/_t o(1) as \quad a \rightarrow o$$

And hence

$$\left| \pi_1(P_t^{(1)} - \pi_1\left(P_t^{(2)}\right) \right| \leq ((r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ P_t^{(1)} - P_t^{(2)} /\!/_t o(t) as \quad t \rightarrow o$$

Which together with a similar bound for $\left| \pi_2(P_t^{(1)} - \pi_2\left(P_t^{(2)}\right) \right|$ leads to

$$\| \pi\,(P_t^{(1)} - \pi\left(P_t^{(2)}\right) /\!/ \leq (r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ P_t^{(1)} - P_t^{(2)} /\!/_{t*} o(t^*)$$

As $t^* \rightarrow o$ showing that if $t^*$ is small enough, the mapping $\pi$ on

$s_{t^*}(\Phi, r^*)$ is a contraction. Hence the system (2.6)-(2,8) has a unique solution for $o \leq t \leq t^* for\ t^*$ sufficiently small and positive

Now if [0,T] is any finite interval on $R^+$, we can consider [0,T] as a finite union of intervals of length less than or equal to $t^*$ Since the estimate (2.11)is valid for any finite interval we can extend the solution of (2,6)-(2.8) form $[0, t^*]$ to [0,T] and this completes the proof,

3. Stationary age distributions and their linear stability we will now establish the existence of time independent solution

$P_t^*, P_t^*$ Of (1.1)-(1.6): such solutions satisfy

(3.1) $\frac{dp\,P_2^*(a)}{da}+f_i(\mathrm{a}, P_1^*, P_2^*)\,P_i^*(a) = 0,\ \ i = 1,2, a > 0,$

(3.2) $P_i^* = \int_o^\infty P_i^*(a)da, i = 1,\ 2,$

(3,3) $P_i^*=\int_o^\infty b_i(a,\ P_1^*,\ P_2^*)P_i^*(a)da,\ i = 1.\ \ 2$

All solutions of (3.1)are of the from

(3.4) $P_i^*(a) = P_i^*(o)\exp\left[-\int_o^\infty f_i(s, P_1^*, P_2^*)ds\right], i = 1,\quad 2, a \geq 0,$

Which together with (3.3) lead to

$1=\int_o^\infty b_1(a, P_1^*, P_2^*)\exp\left[-\int_o^a f_1(s, P_1^*, P_2^*)ds\right]da,$

(3.5) $1=\int_o^\infty b_2(a, P_1^*, P_2^*)\exp\left[-\int_o^a f_2(s, P_1^*, P_2^*)ds\right]da$

Thus the existence of stationary solutions of (1.1)-(1.6) reduces to the existence of a pair of positive constants $P_1^*, P_2^*$ satisfying (3.5). We can now prove the following,

THEOREM 3. Assume that $f_1$ and $b_i$ satisfy ( $A_1$ ) and ( $A_2$ )and let

(3.6) $F_{i.}\ (P_1, P_2) = \int_o^\infty b_i(a, P_1, P_2)\exp\left[-\int_o^\infty f_i(s, P_1, P_2)ds\right]da,$

$i$ =1, 2, $P_1 \geq 0$, $P_2 \geq 0$,

**And suppose the following hold:**

(3.7) $A_4$. $F_{i.}\ (0,\ 0) > 1,\ i = 1, 2$:

$A_5$. *there exists a positive constant c for which*

$F_1$(C,0)< $F_{2.}$(C,0),

(3.8) $F_1$(0.C)< $F_{2.}$(0,C),

Then there exists a unique pair $(P_1^*, P_2^*)$ of real numbers $P_1^* > o$,

$P_2^* > 0$ such that

(3,9) $F_1(P_1^*, P_2^*) = 1 = F_2(P_1^*, P_2^*)$

And the unique nonegative solution of (3.1)-(3.3) is given by

(3.10) $P_2^*(a) = [P_2^* \exp[-f_i(s, P_1^*, P_2^*)ds])/(\int_o^\infty \exp[-\int_o^a f_i(s, P_1^*, P_2^*)ds]da)$,

$i = 1, 2, \quad a \geq 0,$

Proof Consider the elements of the surfaces defined by

$$Z = F_1(P_1, P_2), z = F_2(P_1, P_2), \ P_1 \geq 0, P_2 \geq 0,$$

In the ($P_1$,$P_2$,z) space. By hypothesis ($A_1$) we have $\partial F_i / \partial P_j < 0$ *and* , by ($A_2$), $\partial b_i / \partial P_j \leq 0$ for all

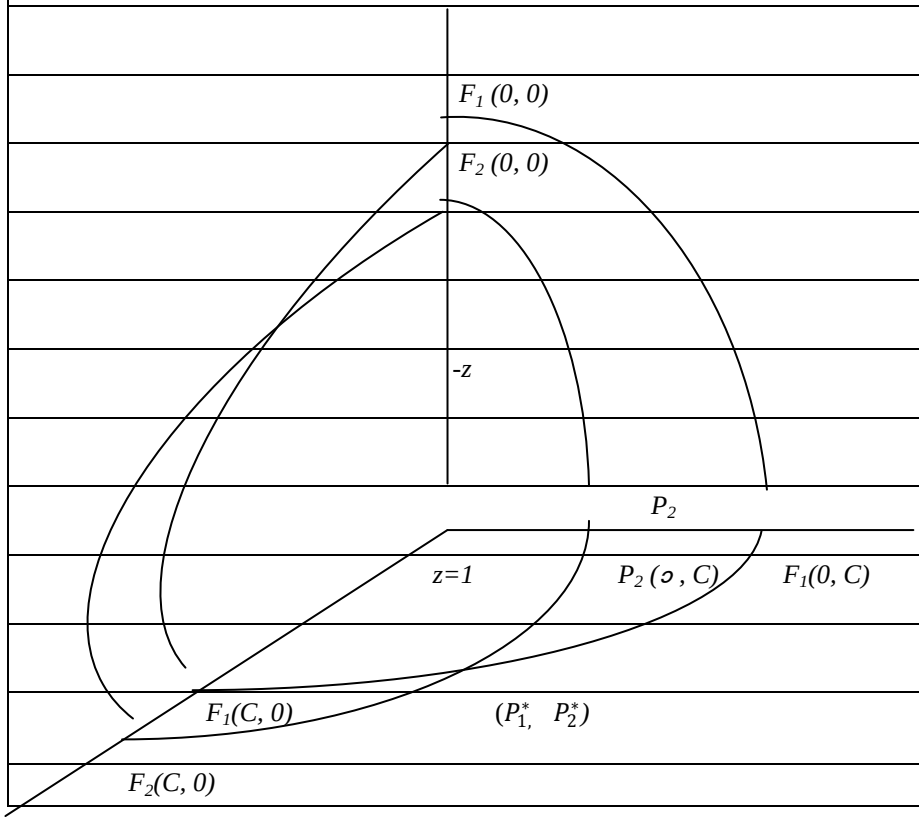
$$P_1 \geq 0, P_2 \geq 0, \quad i = 1, 2,$$

$j = 1, 2,$ Hence by ($A_4$) the intersections of these surfaces with the plane z=1 define two curves on the z=1 plane connecting the I ines {z=1, $P_1 \geq 0, P_2 \geq 0,$} and {z= $1, P_2 \geq 0, P_1 \geq 0$, These two curves are defined by

$$F_1(P_1, P_2 = 1 \ and \ F_2(P_1, P_2) = 1, \ F_1(P_1, P_2 \geq 0,$$

Which by ($A_s$) will intersect at a unique point say $((P_1^*, P_2^*)$ with $(P_1^* > 0, P_2^* > 0$ (see figure). The uniquess of the point $((P_1^*, P_2^*)$ *is a* consequence of the smoothness of the surfaces $z = F_1(P_1, P_2)$ and $z = F_2(P_1, P_2)$

Following is a hypothetical case for Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates attributable or ascribable in the transactions and reactionary potentialities.



(The figure corresponds to the case $F_1(0,0) > F_2(0,0) > 0$. The other possibility $1 < F_1(0,0)$ is treated similarly) It will now

Follow that for such a pair $(p^*_1, p^*_2)$ the unique nonnegative solution of (3.1) - (3.3) is given by (3.10).

Now to examine the linear stability of the stationary age distributions $p^*_1, p^*_2$, we let

(3.11)
$$\begin{cases} p_i(a,t) = p^*_i(a) + u_i(a,t) \\ p_i(t) = p^*_i + p_i(t), \quad i = 1, 2, \quad a > 0, \quad t > 0, \\ p_i(t) = \int_0^\infty u_i(a,t)da, \end{cases}$$

In (1.1)-(1.5) and derive the following variational system after neglecting the nonlinear terms in the perturbations $\boldsymbol{u}_i$ and $\boldsymbol{p}_i$ $(i = 1, 2,)$ (such a procedure of linear stability analysis can be justified As has been done in Gurtin and MacCamy [4].)

(3.12)
$$\begin{cases} Du_i(a,t) = -f_i\left(a, p^*_1, p^*_2\right) u_i(a,t) - p^*_i(a)\left\{\sum_{i=1}^{2} \frac{\partial f_i}{\partial p^*_i} p_i(t)\right\}, \\ u_i(0,t) = \int_0^\infty b_i\left(a, p^*_1, p^*_2\right) u_i(a,t)da + \int_0^\infty p^*_i(a)\left\{\sum_{i=1}^{2} \frac{\partial b_i}{\partial p^*_i} p_i(t)\right\} da \end{cases}$$

Where $i = 1,2$ and (3.13)
$$p_{11}(t) = p_1(t), \quad p_{12}(t) = p_2(t-\tau), \quad p_{21}(t) = p_1(t-\tau), \quad p_{22}(t) = p_2(t)$$

To consider the asymptotic bahaviour as $t \to \infty$ as of solutions of (3.12) we let

(3.14)
$$\begin{cases} u_i(a,t) = u_i(a) \exp[\lambda t], \quad u_i(a) \to 0 \text{ as } a \to \infty, \\ p_i(t) = p^*_i \exp[\lambda t], \\ \qquad\qquad i = 1, 2, \quad a > 0, \quad t > 0, \\ p^*_i = \int_0^\infty u_i(a)da, \end{cases}$$

Using (3.14) in (3.12) - (3.13),

$$\frac{du_i}{da} + \lambda u_i = -f_i\left(a, p^*_1, p^*_2\right) u_i(a) - p^*_i(a)\left\{\sum_{j=1}^{2} \frac{\partial f_i}{\partial p^*_j} p^*_i\right\},$$

(3.15)

$$\mu_i(0) = \int_0^\infty b_i\left(a, p_1^*, p_2^*\right)\mu_i(a)da + \int_0^\infty p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial b_i}{\partial p_j^*}\ p_i^*\right\}da,$$

Where $i = 1,\ 2,$ and

$$p_{11}^* = p_1^*, \quad p_{12}^* = p_2^* \exp\{-\lambda\tau\},$$

$$p_{21}^* = p_1^*,\ \exp\{-\lambda\tau\},\ p_{22}^* = p_2^*.$$

If we integrate the first of (3.15) and use the fact that $\mu_i(a) \to 0$ as $a \to \infty$ we can derive that $p_1^*$ and $p_2^*$ are governed by

(3.16) $$\lambda p_1^* - \int_0^\infty p_1^*(a)\left[\left(\frac{\partial b_1}{\partial p_j^*} - \frac{\partial f_1}{\partial p_1^*}\right)p_1^* + \left(\frac{\partial b_1}{\partial p_2^*} - \frac{\partial f_1}{\partial p_2^*}\right)p_2^* e^{-\lambda\tau}\right]da$$

$$= \int_0^\infty \left[b_1\left(a, p_1^*, p_2^*\right) - f_1\left(a, p_1^*, p_2^*\right)\right]\mu_i(a)da,$$

(3.17) $$\lambda p_2^* - \int_0^\infty p_2(a)\left[\left(\frac{\partial b_2}{\partial p_1^*} - \frac{\partial f_2}{\partial p_1^*}\right)p_1^* e^{-\lambda\tau} + \left(\frac{\partial b_2}{\partial p_2^*} - \frac{\partial f_2}{\partial p_2^*}\right)p_2^*\right]da$$

$$= \int_0^\infty \left[b_2\left(a, p_1^*, p_2^*\right) - f_2\left(a, p_1^*, p_2^*\right)\right]\mu_2(a)da,$$

From the first of (3.15) we get

(3.18) $$\mu_i(a) = \left\{\mu_i(0)p_i^*(a)/p_i^*(0)\exp[\lambda a]\right\} - p_i^*(a)\exp[-\lambda a]\int_0^a \left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right\}e^{\lambda s}ds, \quad i = 1, 2,$$

Which by the second of (3.15) becomes,

(3.19) $$\mu_i(a)\left[p_i^*(a)/p_i^*(0)c_i\exp(\lambda a)\right]\left[\int_0^a p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right\}da\right.$$

$$\left. - \int_0^\infty b_i\left(a, p_1^*, p_2^*\right)p_i^*(a)\,e^{-\lambda a}\left\{\int_0^a\left(\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right)e^{\lambda s}ds\right\}da\right]$$

$$- p_i^*(a)\,e^{-\lambda a}\int_0^a\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right\}e^{\lambda s}ds, i = 1, 2,$$

where

(3.20) $$c_i = 11 - \frac{1}{p_i^*(0)} \int_0^\infty b_i\left(a, p_1^*, p_2^*\right) p_i^*(a) e^{-\lambda a} da, \quad i = 1, 2,$$

Now supplying $\varphi_1(a)$ and $\varphi_2(a)$ from (3.19) in (3.16)-(3.17) we derive that

$$\left(\lambda - A_{11}\right) p_1^* - A_{12} p_2^* \exp\left(-\lambda\tau\right) = g_{11}(\lambda) p_1^* + g_{12}(\lambda) p_2^* \exp(-\lambda\tau),$$

(3.21)

$$-A_{21} p_1^* \exp\left(-\lambda\tau\right) + \left(\lambda - A_{22}\right) p_2^* = g_{21}(\lambda) p_1^* \exp\left(-\lambda\tau\right) + g_{22}(\lambda) p_2^*$$

Where

(3.22) $$A_{i\varphi} = \int_0^\infty p_i^*(a) \left[\frac{\partial b_i}{\partial p_\varphi^*} - \frac{\partial f_i}{\partial p_\varphi^*}\right] da,$$

(3.23) $$g_{i\varphi}(\lambda) = \left[p_i^*(0) c_i\right]^{-1} \left\{\int_0^\infty \sigma_i(a) p_i^*(a) e^{-\lambda a} da\right\} \varphi \left\{\int_0^\infty \left[p_i^*(a) \frac{\partial b_i}{\partial p_\varphi^*} - b_i\left(a, p_1^*, p_2^*\right) p_i^*(a) e^{-\lambda a} \left(\int_0^a \frac{\partial f_i}{\partial p_\varphi^*} e^{\lambda s} ds\right)\right] da\right\}$$

$$- \int_0^\infty \sigma_i(a) p_i^*(a) e^{-\lambda a} \left(\int_0^a \frac{\partial f_i}{\partial p_\varphi^*} e^{\lambda s} ds\right) da, \quad i, \varphi = 1, 2,$$

(3.24) $$\sigma_i(a) = b_i\left(a, p_1^*, p_2^*\right) - f_i\left(a, p_1^*, p_2^*\right), \quad i, = 1, 2.$$

To solve the linear perturbational system (3.12) it is enough to find the constants $\lambda, p_1^*, p_2^*$ satisfying (3.21) since one can then use (3.19) to find the perturbations from (3.14). A nontrivial solution $\left(p_1^*, p_2^*\right)$ for (3.21) will exist if and only if $\lambda$ is a root of the equation $et \begin{bmatrix} \lambda - \left(A_{11} + g_{11}(\lambda)\right) & -\left(A_{12} + g_{12}\right) e^{-\lambda\tau} \\ -\left(A_{21} + g_{21}\right) e^{-\lambda\tau} & \lambda - \left(A_{22} + g_{22}(\lambda)\right) \end{bmatrix} = 0$

Or equivalently

(3.25) $$\lambda^2 - \lambda\left(A_{11} + A_{22}\right) + A_{11}A_{22} - A_{12}A_{21} \exp\left[-2\lambda\tau\right] = \lambda s_1(\lambda) + s_2(\lambda, \tau)$$

Where $S_1(\lambda) = g_{11}(\lambda) + g_{22}(\lambda)$ ,

$S_2(\lambda, \tau) = \left[A_{21} g_{12}(\lambda) + A_{12} g_{21}(\lambda) + g_{12}(\lambda) g_{21}(\lambda)\right] \exp\left[-2\lambda\tau\right]$

$$-A_{11} g_{22}(\lambda) - A_{22} g_{11}(\lambda) - g_{11}(\lambda) g_{22}(\lambda)$$

The following result establishes the linear local asymptotic Stability of the stationary age distributions $p_1^*(a)$ and $p_2^*(a)$

THEOREM 4 Assume the following:

$A_6.\ 0 < \beta_i^* < d_i^*\ (i,\ =1,\ 2.)\ \ (see(A_1)\ and\ (A_2));$

$A_7.\ \{A_{11} + A_{22} + \| S_1(0) \|\}^2 < 4\{(A_{11}A_{22} - A_{12}A_{21}) - \| S_2(0) \|\},$

$$\| S_1(0) \| = |g_{11}(0)| + |g_{22}(0)|,$$

$$\| S_2(0) \| = |A_{11}\, g_{12}(0)| + |A_{12}\, g_{21}(0)| + |g_{12}(0)\, g_{21}(0)|$$

$$+ |A_{11}\, g_{22}(0)| + |A_{22}\, g_{11}(0)| + |g_{11}(0)\, g_{22}(0)|$$

Then all the roots of (3.25) have negative real parts and hence the

perturbations $u_i\ (a,\ t) = \phi(a) \exp[\lambda t] \to 0\ as\ \ t \to \infty$ for $i = 1, 2\ and\ a \geq 0$

Proof . Define $F_1(\lambda, \tau)$ and $F_2(\lambda, \tau)$ as follows:

$F_1(\lambda, \tau) = \lambda^2 - \lambda(A_{11} + A_{22}) + A_{11}A_{22} - A_{12}A_{21} \exp[-2\lambda\tau];$

$F_2(\lambda, \tau) = \lambda S_1(\lambda) + S_2(\lambda, \tau),$

Using $(A_6)$ one can show that for $\lambda$ with re $(\lambda) \geq 0$, $c_1$ and $c_2$ are positive and bounded away from zero, Now since $A_{ij} > 0$ and $\tau \geq 0$ we find from the nature of the dependence of $g_{ij}$ on $\lambda$ that for all $\lambda\ \ \mu \pm iw$ with $\mu \geq 0$ and $w \geq 0$ we have

(3.26) $$\left| F_1(\lambda, \tau) - P_2(\lambda, \tau) \right|$$

$$\geq \left| F_1(\lambda, 0) | - F_2(\lambda, \tau) \right|$$

$$\geq |\lambda|^2 - |\lambda|\{A_{11} + A_{22} + \| S_1(0) \|\} + \{A_{11}A_{22} - A_{12}A_{21} - \| S_2(0, \tau) \|\}$$

$$\geq \left[(|\lambda| - P)^2 + (q)\right]^2$$

Where

$p = \{A_{11} + A_{22} + \| S_1(0) \|\}/2,$

$q = \left\{4 | (A_{11}A_{22} - A_{12}A_{21}) - \| S_2(0, T) \| | - | A_{11} + A_{22} + \| S_1(0) \| |^2\right\}/2,$

It will now follow, from (A) and (3.26), $F_1(\lambda, \tau) = F_2(\lambda, \tau)$ connot have roots with zero or positive real parts and this completes the proof.

In conclusion we remark that the condition (3.7) means that the net creation rates of each system is greater than utility; a condition of this type is well known in age wise classification of particles, atoms, molecules, cosmological phenomenon. (3.8) can be interpreted to mean that the intraspecific inhibition is higher than that of the

interspecific interaction. The conditions of Theorem 4 are analytical and do not lend themselves for any worthwhile interpretation in terms of the model parameters. So also an interpretation can be recorded as a remark in respect of destruction rates of particles, properties, elements, or projections.

Note: The model is applicable to each and every module. As stated in the foregoing, there are possibilities of situation where an account itself be a receiver and it is also the one to which debit is raises. Such examples are legion in all the fields of finance, economics, political science and physics and chemistry, and of course mathematics. Classification based on time is the preeminently and predominantly suited premise for such model. We repeat that the model is to be read with each module We will expatiate the importance and implicational ramifications as we proceed with the work. Assortment-blend, collection-combinations, combo-gallimaufry, goulash-hash, hodgepodge-medley, mishmash-mixture, motley-pastiche patchwork-salmagundi, soup-stew, mixed bag-mélange, olio- olio podrida, assortment variety, Array choice, collection combination, combo diversity, hodgepodge jumble, kind medley, miscellany mishmash, mixed bag mixture, mélange potpourri, selection sort, brew concoction, beverage blend, broth compound, distillation drink, fermentation hash, hodgepodge infusion, instillation liquor, mélange miscellany, mishmash mixture, mélange potpourri, borscht bouillon, bowl brew, chowder concoction, consommé decoction, gumbo hodge-podge, olio porridge, potage potpourri, pottage puree, splash stock, vichyssoise hodgepodge, mixture- mess, collection combination, goulash hash, jumble medley, miscellany mishmash, mixed bag mélange, olio patchwork, potpourri salmagundi, jumble clutter, disarrangement disarray, disorder farrago, gallimaufry garbage, goulash hash, litter medley, mess miscellany, mishmash mixture, muddle mélange, olio pastiche, patchwork potpourri, salmagundi scramble, shuffle snarl, tangle tumble, medley miscellany, assortment brew, collection combo, conglomeration farrago, hodgepodge jumble, melee mingling, mishmash mixed bag, mixture mélange, pasticcio pastiche, patchwork potpourri, salmagundi variety, mélange mixture, assortment combo, farrago gallimaufry, pasticcio pastiche, patchwork potpourri, salmagundi soup, stew miscellany, accumulation aggregation, anthology assortment, brew collectanea, combination combo, compilation conglomeration, accumulation diversity, farrago gallimaufry, melee mishmash, mix mixed bag, mixture muddle, mélange odds and ends, olio pasticcio, pastiche patchwork, potpourri salad, salmagundi smorgasbord, stew variety, mishmash jumble, medley mess, miscellany mixed bag, mixture mélange, olio patchwork, potpourri salmagundi, mixed bag array, assortment choice, collection combination, combo conglomerate, diversity hodgepodge, jumble medley, mélange miscellany, mishmash mixture, montage mélange, odds and ends, patchwork potpourri, smorgasbord sundries, mix/mixture assortment, combination, admixture adulteration, alloy amalgam, amalgamation assimilation, association batter, blend brew, combine combo, commixture composite, compound concoction, confection conglomeration, cross crossing dough fusion, goulash grab bag, hodgepodge hybrid, hybridization incorporation, infiltration interfusion, jumble mash, medley merger, mingling miscellany, mishmash mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, odds and ends assortment, bits and pieces, debris etcetera, hodgepodge jumble, knickknacks leavings, leftovers litter, medley melee, miscellaneous paraphernalia, miscellany mishmash, mixed bag motley, mélange notions, novelties oddments, olio particles, patchwork potpourri, rags refuse, remainder remains, remnants rest, rummage scraps, sundries sundry items, compilation mishmash, paste-up patchwork, potpourri reappropriation, reproduction synthesis, conglomeration cross, crossing dough, fusion goulash, interfusion jumble, mash medley, merger mingling, miscellany mishmash, mix mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, colluviums gallimaufry, hash mixture, pasticcio pastiche, conglomeration cross section, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multiformity multiplicity, mélange potpourri, range shift, variegation variousness, divergencies diverseness, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multifariousness multiformity,

mélange potpourri, range shift, soup stew, variance variation, variegation variousness, oddment odds and ends, assortment bits, polymorphism assortment, change collection, combo conglomeration, cross section departure, discrepancy disparateness, divergencies diverseness, diversification diversity, fluctuation heterogeneity, heterogeneousness incongruity, intermixture many-sidedness, medley miscellaneousness, miscellany mishmash, mixed bag mixture, notions novelties, oddment oddments, olio particles, patchwork potpourri of the particles, interactions, accounts, transactions in the cosmos. The wisdom and understanding of our own divine essence is like a beaconing begorra, bureaucratic burdock, befuddling binary and digital digamma, that invites us inward toward ever deeper connection with all around us. Yet rather than look within, most in this world prefer to look outward for divine wisdom, thus sustaining a dependence upon a vast hierarchy that stretches between the individual and the Divine. In all our wanderings away from the Divine, humankind has obscured, obfuscated, obliterated, adjured, abnegated and abjured its most compelling features through a persistent belief systems in limitations arising out of the controlling mechanisms of this hierarchy of sub ordinates, superordinates and coordinates. Has my consciousness changed? Why that doubt? Twist and shake, ponderous novelty themes, pathos, poignancy in writing in different styles, accord-coetaneousness, coevality-coincidence, concurrence-conformity, conjunction-contemporaneousness, harmony, order, peace, simultaneousness- synchronicity exciting optimism of youth, power obmbudsman,pulsations and throbbing of pop, hard luck stories of people, modernity and mischief, grimalkins of rurality, wobbling boats of lovers, saxophone in mouths all point to newer value orientations. The Divine dances and the “Grand design “is perpetrated outside of the confines of any hierarchical structures, be it feudal, organizational. It is complete within itself, and has a singular purpose of demonstrating the collective potential of all life within the universe. It is the archetype of perfection. It is the standard bearer of each soul's innate design and ultimate destiny. The essence of the Divine is far beyond mental conception, yet humanity's tendency is to resort to the limiting language of the hierarchical paradigm acceding- accession, accommodation- accord, accordance, acknowledging, adjustment-affiliation, affinity-alliance, amity-approving, arbitration-arrangement, assenting, authorizing, bargaining-compatibility, compliance--compromise, concert-concession, conformity-congruity, consistency-correspondence, endorsing, granting-harmony in conformity to pen chance ,predilections of the protagonist to define and understand it. Searching for Connection and Wholeness calls for not lumping elements like me made of clay feet and weak knees, owing allegiance to the fringe, without courante and courage of conviction and will for vindication, visual restlessness and exaggerated narratives of self to its own, exciting potentialities available, terrestrial passion, and skew of scams all bear an infallible demonstration for new order of consciousness. Theorists have posited a virtual smorgasbord of objects to account for dark matter, although many of them have fallen prey to observational constraints. As leading possible candidates for baryonic dark matter, there are black holes (large and small), brown dwarfs (stars too cold and faint to radiate), sun-size MACHOs, cold gas, dark galaxies and dark clusters, to name only a few. The range of particles that could constitute nonbaryonic dark matter is limited only slightly by theorists' imaginations. The particles include photinos, neutrinos, gravitinos, axions and magnetic monopoles, among many others. Of these, researchers have detected only neutrinos--and whether neutrinos have any mass remains unknown. Experiments are under way to detect other exotic particles. If they exist, and if one has a mass in the correct range, then that particle might pervade the universe and constitute dark matter. But these are very large "ifs." To a great extent, the details of the evolution of galaxies and clusters depend on properties of dark matter. Without knowing those properties, it is difficult to explain how galaxies evolved into the structures observed today. As knowledge of the early universe deepens, I remain optimistic that we will soon know much more about both galaxy formation and dark matter. What we fail to see with our eyes, or detectors, we can occasionally see with our minds, aided by computer graphics. Computers now play a key role in the search for dark matter. Historically, astronomers have focused on observations; now the field has evolved into an experimental science. Today's astronomical experimenters sit neither at lab benches nor at telescopes but at computer terminals. They scrutinize cosmic simulations in which tens of thousands of points, representing stars, gas and dark matter; interact gravitationally over a galaxy's lifetime. A cosmologist can tweak a simulation by adjusting the parameters of dark matter and then watch what happens as virtual galaxies evolve in isolation or in a more realistic, crowded universe.

Old consciousness will not do. Newer one has to be imbibed and cultivated. When people are unaware of their inner wisdom and wholeness, they tend to search for order and security outside themselves. That is my personal experience. There is always search for sanctity of emotion, aesthetic expression, pursuit of harmony, unraveling realms, syncopated ending and stillness of mind and shiver of heart at the mirror broken somewhere are the strains of newer wisdom if it can be called one. Same questions shall be asked and same shall be answer in the similar vein. Uncertain of their place within the hierarchical order of the world, they define themselves based upon their insecurities. Individuals thus become only pieces of their wholeness and like shards of glass from a beautiful vase; they bear little resemblance to their aggregate beauty. Within the hierarchy, many in high positions of power have taken advantage of our collective insecurities in attempting to guide the development of all humankind. They have obscured the direct connection between the individual and the Divine through a variety of means designed to intercede between our inner essence and our divine source. Each individual can come to know them to be free of all forms of institutional and organisational and societal hierarchical control. This is not to imply that we should not trust others or join together in bonds of friendship and community. It is simply a reminder that relative truth is constantly shifting in the hands of those who desire to control. And even when the motive for their controlling behaviors may be of good will, it is still a form of control. When the revealers of "truth" within the hierarchy withhold and suppress information, they are usually positioning themselves to maintain or acquire more power rather than to disseminate empowerment to all. The desire for connection and wholeness is a fuel that drives us to seek out and explore the hierarchical paradigm. This inner longing provides us with the motivation to seek help and guidance from a specific group within the hierarchy, and in so doing, cultivate a sense of belonging and connection. Furthermore, the hierarchical paradigm is a stage whereby we develop a sense of connection to some grand, encompassing vision. This is why the hierarchy nurtures prophets who point toward a greater vision. Spiritual leaders are able to peer deeply beneath the surface reality of life and experience how intricately connected every life form is, and how the composite of all life is intelligent far beyond measure. These visionary leaders can thus interpret reality through their personal abilities to perceive and express life's dimensional depth and limitless intelligence. Yet no one is able to articulate life's full dimensional depth and breadth with the tools of language. They can only, at best, describe their interpretation or their impressions. In actuality, all of us are able to peer beneath the surface reality of life and perceive a unique vision of the universe. We require only time and intention to develop our own interpretations. And this is precisely what many great spiritual leaders have taught. Life's deeper essence is not an absolute to be experienced by the chosen few, but an evolving, dynamic intelligence that wears as many faces as there are life forms. No group or people have the exclusive portal into the universe by which the Divine expresses itself in all its majesty. The portal is open and available to all, because the Divine is within all things. Those recognized as great prophets each produced a vision of the universe beyond what was currently defined by the hierarchy. Because their interpretations were articulated with authority and depth of insight, they became a target of debate among various groups in the hierarchy. This debate then created a polarity of belief. A sympathetic constituency emerged to defend and embellish their leader's interpretation, while established groups held it in contempt of previously held beliefs. Invariably, the leader's vision became confined and shaped into dogma by followers who desired to create a new religion or sect. Thus, this infusion of fresh insight quietly receded into the hands of the hierarchy, where its deeper meaning was obscured by the very fact that it was incorporated into a rigid structure that both protected and promoted it. A new paradigm is emerging that promotes a clear connection of individual consciousness to the compelling features of the Divine without the intervention of a hierarchy. This is when the fables and myths of history step into the light and become known as they were originally intended. This is the time when language will be transformed into a new form of communication that breaks down all barriers of control. Personal transformation, through the awakening of inner wisdom and understanding, is the pathway into wholeness. The transformation paradigm is initiated simply by the recognition that rather than the dependency-inducing ways of the hierarchical paradigm, there are accelerated, independent pathways that bypass the hierarchy and lead to self-mastery. These new pathways lead to the divine wisdom and understanding that is present within all of us. This wisdom can be accessed through the practice of three principles of transformation: seeing the Divine in

all, nurturance of life, and gratitude. The application of these life principles disengages individuals from the controlling elements of the hierarchy, thus initiating the transformation experience. As there are relative truths, there are relative freedoms. As individuals evolve through the hierarchical process, an ever-increasing sense of freedom is gained, yet external forces continue to exert control through limiting language, confining belief systems, and more. These controlling influences lead to anatomical architecture, body, build, cagey carcass, construction, enclosure, fabric, flounce, form, framework, fringe, groundwork, hem, mount, mounting outline, physique, scaffold, scaffolding scheme, setting shell, stage, structure, support, system, trim, trimming, truss- valance continued reliance upon the hierarchy as it unceasingly attempts to impose a sense of inequality between us. The underlying equation of the hierarchical paradigm is: individual + hierarchy = God connection. In the case of personal transformation it is: individual + inner wisdom and understanding divine equality with all. The time has come to integrate the dominant paradigm of the hierarchy with the liberating transformation paradigm. That shall redeem all promises made to oneself. There shall be a Matrix force or balancing force that shall remain UN satisfies until conservativeness is reached. This integration occurs naturally once we have fully explored the two paradigms and develop a synthesis paradigm whereby transformation is attained by searching for connection and wholeness through our own inner wisdom and understanding. It is this combination of release from dependency on the hierarchy and transformation from within that initiates the synthesis paradigm. Once we take responsibility for our own transformation and integration, it does not mean that the hierarchy is to be shunned or avoided. The hierarchy is quite benign as a manipulative force and merely represents one important stage in the journey toward wholeness. What is being set in motion now is the initial preparation for these paradigm shifts. More specifically, these paradigms will be simultaneously played out over the coming years. As always, it will be the choice of each one of us as to which paradigm we embrace in our journey. All of the highest imaginings of humanity are yet unaware of our deepest foundation. We have sought the upper reaches of the building, yet remain unaware of the foundation's design. It is here, at the very core of existence that the Divine is bursting forth with its creative energy while simultaneously reintegrating with its invitation to wholeness. It is here that equality is realized, not in the lofty places of relative truth lodged in the hierarchy, but rather in the deepest part of the foundational plan of life's original source and ultimate destiny. The origin and destiny of existence is the tone of equality in life. Listen for this tone – this frequency of vibration – and follow it back into the very foundation from whence all things arise and return. Consider these words as symbols only. Feel the truth that stands behind these words, and tap into this empowering energy force that reaches out for you. Know it as a tone or vibration – a resonance that waits for you around every corner in which your life will turn. It is a beacon of the Divine gathering itself into the form of language in order to usher you to a place from which you can experience the formless tone of equality – the bypass of limitation. It is the primal language of our divine source that bestows to you the freedom to generate your own deepest beauty in the expression of your highest truth.(Reference: Wikipedia)

## Gopalasamy Model For Procrastinated Conducive, Felicitous Counter Sallies In Two Specified Systems

Abstract

Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of decay rate of perturbations. Sole intention, main objective and primary concern of

the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are counted in the same vein as that of predator prey. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 1. Preliminary remarks

Almost all the equations of population ecology have been derived starting from the following simple format (Note this can be used for the development of the creed and type of systemic types in case of cosmology and astrophysics and fiancé and econometrics. To give an example of finance, we can think of a account of the company giving rise to various accounts of its employees, and that is the contribution of the single account for the development of the total system types of the given configuration). It may be account, particle, atom, molecule, proposition, problem or any such thing. We shall delineate and disseminate information and particulars as and when the model shall be applied to the problem in question.

$$\frac{dN(t)}{dt} = \text{an individual's contribution}$$

*to population growth* $N(t)$

(1.1)

where $N(t)$ denotes the size of a population of a single species (or biomass) with overlapping generations at time $t$, and $dN(t)/dt$ is the rate of change of the population size at time $t$. Subsequently one makes an assumption regarding the factor in the brackets in (1.1); in particular if one assumes that an individual entity's contribution to the population growth is denoted by a function say, f(t, N), defined suitably for all t>0, N then one obtains from (1.1) the so called Kolmogorov formulation in the form

$$\frac{dN(t)}{dt} = f(t, N(t))\, N(t).$$

(1.2)

Various choices of $f$ together with some ecologically plausible assumptions such as the temporal constancy of the environment and density dependent effects in $f$ lead to various equations. For instance if $f(t, N) \equiv r$ (a positive constant) one obtains the Malthusian formulation in the form

$$\frac{dN(t)}{dt} = rN(t)$$

And if one assumes $f(t, N) \equiv r - (r/K) N$ for some positive constants r and K one gets the familiar logistic equation of population biology

$$\frac{dN(t)}{dt} = rN(t) \left\{ 1 - \frac{N(t)}{K} \right\}$$

In order to device models with time delayed effects, such effects are introduced in the average growth rate $F(t, N)$ of (1.2); for instance if $f(t, N) \equiv r - (r/K) N(t - \tau)$ for some constant $\tau > 0$ we obtain Hutchinson's[19]delay loistic equation

$$\left\{ \frac{dN(t)}{dt} = rN(t) \left\{ 1 - \frac{N(t-\tau)}{K} \right\} \right. .$$

(1.3)

If for instance

$$F(t, N) = r - (r/K) N(t) - \int_{-\infty}^{t} k(t - s) F(N(s))\, ds$$

Where $k$ is a piecewise continuous nonnegative valued function such that $\int_0^\infty k(s)\,ds = 1$ with some appropriate function $F$, we will get Volterra's model of a population which pollutes its environment and pollution has an accumulative and toxic effect; such an equation is of the form

$$\frac{dN(t)}{dt} = N(t) = \left\{ r - (r/K) N(t) - \int_{-\infty}^{t} k(t-s) F(N(s))\,ds \right\} \quad (1.4)$$

Which has been studied in different formats by several authors (Cushing [12], Cohen *et al.* [10], and Gopalasamy [16])

One of the significant advantages of deriving various models from the prototype (1.2) is that if the initial; conditions (initial population sizes) are nonnegative then the nonnegativity of the population size $N(t)$ for $t > 0$ follows from the fact that solutions of (1.2) satisfy

$$N(t) = N(0) \exp\left[\int_0^t f(s, N(s))\,ds\right]$$

Which implies that $N(t) \geq 0$ if $N(0) \geq 0$.

While the time delayed logistic equation (1.3) explains some of the observed oscillations of certain populations, its derivation has not been entirely satisfactory. For instance it has been pointed out by Ricklefs ([27], page 488) that "time lags might be expected to occur primarily in stabilizing processes involving production rather than death because death is an immediate response to environmental change". Recently Cushing [13] has examined the effects of response and maturation delays in a number of single and multi system models; as it has been noted by Cushing [13], the models involving delayed reproduction or not necessarily derivable from the Kolmogorov type systems of the form (1.2) by suitably defining the average growth rate $(1/N)(dN/dt) = f(t, N(t))$. While Cushing [13] has been mostly concerned with the possibility of delay induced oscillations we will consider a class of delays in reproduction and interspecific interaction which maintain the stability of the systems.

Our starting point to drive models with delayed reproduction is the following balance equation which assumes that there is no immigration or emigration. Account, particle, atom, molecule, proposition, problem et al

$$\frac{dN(t)}{dt} = \text{birth rate} - \text{death rate}. \quad (1.5)$$

Now for instance if we consider a population of adult flies then the birth rate at time $t$ depends on the adult population at time $t^{-}\ \tau$ where $\tau$ is the time required for the larvae to become adult flies. If the birth and death rates are governed by density dependent factors then we have from (1.5) the following time delayed model

$$\frac{dN(t)}{dt} = b\big(N(t^{-}\tau)\big) - m(N(t))$$

(1.6)

Where the functions $b$ and $m$ denote the density dependent birth and death rates of Account, particle, atom, molecule, proposition, problem et al respectively. If the time delay is continuously distributed then one can consider in the place of (1.6) an equation of the form

$$\frac{dN(t)}{dt} = b\int_{-\infty}^{t} k\{(t-s)N(s)ds - m(N(t))$$

Or equivalently

$$\left\{\frac{dN(t)}{dt} = b\int_{0}^{\infty} k(s)\right\}N(t-s)ds - m(N(t))$$

(1.7)

Where the delay kernel $k$ denotes the weight of the delayed of the delayed effects.

It is not always obvious that suitable nonnegative initial conditions for (1.6) and (1.7) will imply the nonnegativity of the solutions of (1.6) and (1.7). The purpose of this article is to show that in two species systems of competitive and prey- predator communities, delayed reproduction and delayed interactions donot destabilise the systems if the self- regulating negative feedback effects are sufficiently strong compared to interspecific interactive effects and if the self regulating effects are realised with no time delays. When we say that time delays do not destabilise, we will mean that delays do not render an asymptotically (locally) stable steady state unstable in such a manner that the loss of stability either leads

To a delay – induced bifurcation to persistent and undamped oscillations (as in the case of Hopf- bifurcation) or the relevant linear variational system has unbounded solutions; that is a mathematical instability is not induced by delays in the sense, that the roots of the characteristic equation are pushed to the imaginary axis or to the right of it on the complex plane. However delays can make an otherwise stable system less stable which means that the rate of convergence to zero of solutions of the variational system can decrease with an increase of delay and the variational system can even oscillate (perhaps violently) before converging to the trivial state. Such a reduction in stability can

happen even when all the roots of the characteristic equation have negative real parts. For our purposes such a reduced stability is still stability. In order to see the effects of delays on such less stable systems we will propose below a method of estimating the rate of decay of local perturbations based on the variational system and indirectly relying on the characteristic equation. For a special class of two species models with a single delay we have obtained (Gopalsamy [16]) estimates of decay rates (or return times) usingthe characteristic equation; such a procedure is almost impossible in the case of many species systems with several delays. The method we propose in Section 6 below for the estimate of decay rates can be adapted to *n*- species systems with $n^2$ discrete delays.

Let $x_1(t)$ and $x_2(t)$ denote the population sizes (or biomasses) of two systems competing for a common pool of resources.. Let $b_i$ and $m_i$ $(i = 1, 2)$ denote the respective density dependent birth (potential) and death rates so that in the absence of any time delays the population sizes are governed by (of Account, particle, atom, molecule, proposition, problem et al). In addition as usual, we assume that the system studied follow Haimovici's equations, for which the dove tailing explanation has been appended

$$\left.\begin{aligned}\frac{dx_1(t)}{dt} &= b_1\left(x_1(t)\right) - m_1(x_1(t)), x_2(t))\\ \frac{dx_2(t)}{dt} &= b_2\left(x_2(t)\right) - m_2(x_1(t)), x_2(t))\end{aligned}\right\} \quad (2.1)$$

In order to make the system (2.1) represent a competition system we will make the following assumptions on the birth and death rates; $b_1$, $m_1$ $(i = 1, 2)$ are continuous with continuous partial derivatives for all $x_1 \geq 0$ $(I = 1, 2)$; also we assume (see for instance Rescigno and Richardson [26]) $\frac{\partial b_i}{\partial x_1} > 0, \frac{\partial m_1}{\partial x_j} > 0$ for $x_i > 0$ $(i, j = 1, 2)$, (2.2)

(i) $b_1(0) = 0$, $m_1(0, x_2) \equiv 0, b_2(0) = 0$, $m_2(x_1, 0) \equiv 0$, (2.3)For some $x_{1,}^* > 0, x_{2,}^* > 0$ we have

$$\left.\begin{aligned}b_1(x_1^*, -m_1(x_1^*, 0) &= 0,\\ b_2(x_2^*, -m_1(0, x_2^*) &= 0,\end{aligned}\right\} \quad (2.4)$$

(ii) There exist positive constants $\delta_1$, $\delta_2$ such that

$b_1(\delta_1) - m_1(\delta_1, x_2) < 0$,

$b_2(\delta_2) - m_2(x_1, \delta_2) < 0$, for $x_1 \geq 0$, $x_2 \geq 0$,

(2.5)

(iii) There exists a point $(\alpha, \beta), \alpha > 0, \beta > 0$ such that

$$b_1(\alpha) - m_1(\alpha, \beta) = 0 = b_2(\beta) - m_2(\alpha, \beta).$$

(2.6)

The conditions on $b_i$ in (2.2) mean that the potential (not average) birth rates are density dependent and any negative density dependent feedbacks to the opened or birth rate of Account, particle, atom, molecule, proposition, problem et al rates are included in the destruction rates; the conditions on $m_i$ indicate intraspecific and interspecific confederation needs and necessities. (2.3) implies that (0, 0) is a (trivial) steady state of the system (2.1) as it is in the case of all ecological models; (2.4) means that in the absence of either of the two species the other one has a nontrivial steady state; (2.5) implies that each of the two populations cannot grow unboundedly since when $x_1(t) > \delta_i, \frac{dx_1}{dt} < 0$. (2.6) means that there exists a nontrivial steady state $(\alpha, \beta)$ in the interior of the positive quadrant of the $(x_1, x_2)$- plane. These assumptions are standard and are minimal in the opinion of the author.

An immediate example of (2.1) –(2.6) is the familiar Volterra –Lotka competition system described by

$$\frac{dx_1(t)}{dt} = r_1 x_1(t) - a_{11} x_1^2(t) - a_{12} x_1(t) x_2(t),$$

$$\frac{dx_2(t)}{dt} = r_2 x_2(t) - a_{21} x_1(t) x_2(t) - a_{22} x_2^2(t),$$

(2.7)

Where $r, a_{ij}, (i, j = 1,2)$ are positive constants satisfying the condition

$$\frac{a_{11}}{a_{21}} > \frac{r_1}{r_2} > \frac{a_{12}}{a_{22}}.$$

One of the first questions to be resolved for (2.1) is the following: if $x_i(0) > 0$, $x_2(0) > 0$ does it follow that the corresponding solutions of (2.1) satisfy the condition $x_1(t) \geq 0$, $x_2(t) \geq 0$ for those $t$ for which a solution of (2.1) exists? In the $(x_1, x_2)$ phase plane of (2.1) we have (0, 0), $(x_1^*, 0)$, $(0, x_2^*)$, $(\alpha, \beta)$ as the steady states of (2.1) and also it will follow from (2.3) that the co ordinate axes

$x_1 = 0$ and $x_2 = 0$ are invariant sets f (2.1) ; such invariance of the coordinate axes will imply that if $x_1(0) > 0$, $x_2(0) > 0$ then the corresponding population trajectory $(x_1(t), x_2(t))$ of (2.1) cannot reach the outside of the nonnegative quadrant $x_1 \geq 0$, $x_2 \geq 0$.

Our plan is to derive first a set of sufficient conditions for the coexistence of the two competing species and then in the next section show that under the same conditions the two species can coexist if there are time delays in the reproduction and interspecific interactions. We recall that the two competing species described by (2.1) are said to coexist if and only if the nontrivial steady state $(\alpha, \beta)$ is (at least locally) asymptotically stable. To examine such a stability of $(\alpha, \beta)$ we derive the corresponding linear variational system for the perturbations $X_1, X_2$ where

$$x_1(t) \equiv \alpha + X_1(t), \qquad x_2(t) \equiv \beta + X_2(t)$$

Such a variational system is found to be

$$\left.\begin{aligned} \frac{dX_1(t)}{dt} &= \beta_{11} X_1(t) + \beta_{12} X_2(t), \\ \frac{dX_2(t)}{dt} &= \beta_{21} X_1(t) + \beta_{22} X_2(t), \end{aligned}\right\} \qquad (2.8)$$

Where

$$\beta_{1J} = \left(\frac{\partial b_1}{\partial x_J} - \frac{\partial m_1}{\partial x_j}\right) \text{ evaluated at } (\alpha, \beta) \qquad (i, j = 1,2).$$

(2.9)

The steady state $(\alpha, \beta)$ of (2.1) is locally asymptotically stable if and only if the trivial solution (0, 0) of (2.8) is asymptotically stable. We can derive the following result from elementary considerations.

In the system (2.1)-(2.6) assume the following:

$$\left.\begin{aligned} \frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1}, \\ \frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1}, \end{aligned}\right\} \qquad \text{at } (\alpha, \beta)$$

(2.10)

Then the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (2.1) – (2.6).

The proof follows easily if we note that the characteristic equation associated with the linear system (2.8) is given by

Det $\begin{pmatrix} \lambda-\beta_{11} & -\beta_{12} \\ -\beta_{21} & \lambda-\beta_{22} \end{pmatrix} = 0$

Or equivalently,

$$\lambda^2 - \lambda\,(\beta_{11} + \beta_{22}) + \beta_{11}\beta_{22} - \beta_{12}\beta_{21} = 0. \qquad (2.11)$$

It follows from (2.2) and (2.10) that,

$$\beta_{11} < 0, \beta_{11}\beta_{22} - \beta_{12}\beta_{21} > 0$$

And hence the roots of (2.11) have negative real parts implying the asymptotic stability of the trivial solution of (2.8)

If apply the condition (2.10) to the Volterra-Lotka model (2.7) we note that when a nontrivial steady state say $(\alpha, \beta)$ exists for (2.7), (2.10) leads to

$$a_{11}\alpha > a_{21}\,\beta \; and \; a_{22}\,\beta > a_{12}\,\alpha,$$

Which together imply $a_{11}\,a_{22} > a_{12}a_{21}$. If $(\alpha, \beta)$ exists it is known that $a_{11}\,a_{22} > a_{12}a_{21} > 0$ is a sufficient condition for the (local) asymptotic stability of $(\alpha, \beta)$.

The condition (2.10) which guarantees the coexistence of competing populations in (2.1)-(2.6) simply means that the intraspecific self – regulating feedback effects $\partial m_i / \partial x_1\ (i = 1, 2)$ bare higher than its own positive feedback $\partial b_1 / \partial x_1\ (i = 1, 2)$ together with its negative (competitive) effects on its competitor $\partial m_I / \partial x_J\ (i,\ j = 1, 2; i \neq j)$.

## 2. Coexistence under delayed production and interactions, transactions

We will now establish that two competing species with a sufficiently strong intraspecific regulation can coexist in the presence of delayed reproduction and delayed interspecific interactions. It is well known that if there are delays in the self – regulating intraspecific effects, stationary coexistence may not always be possible (Gopalsamy [16], Gopalsamy and aggarwala [15], Shibata and Saito [28]). Delayed reproduction is more general and widespread in most biological populations. To build a competition model with such delays we proceed as follows.

Let $\tau_{ij}\ (i, j = 1, 2)$ be a set of nonnegative constants with $\tau = \max\ \tau_{ij}\ , i, j = 1, 2,$ and suppose that the two ompeting species display delayed reproduction and delayed interspecific interactions with no delays in the self – regulating negative feedback effects. Such a competition system in a constant environment can be modeled by the autonomous delay differential system

$$\frac{dx_1(t)}{dt} = b_1\big(x_1(t-\tau_{11})\big) - m_1\big(x_1(t), x_2(t-\tau_{12})\big),$$

$$\frac{dx_2(t)}{dt} = b_2\big(x_2(t-\tau_{22})\big) - m\big(x_1(t-\tau_{21}), x_2(t)\big), \qquad (3.1)$$

Where both the creation and destruction rates and the morality rates $b_1$ and $m_1\ (i = 1, 2)$ satisfy the same conditions (2.2)-(2.6) as in the previous section. Along with (3.1) we will

Suppose that the initial population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) sizes are provided in the form

$$x_i(s) = \phi_i(s) \geq 0, \qquad s \in [-\tau, 0], \qquad i = 1,2. \tag{3.2}$$

(Note that the specification of $\phi_i$ need not be on the entire interval$[-\tau, 0]$; but such a hypothesis is convenient.) Since (3.1) is not of the standard Kolmogorov-type, we have to establish that the solutions of (3.1) will remain nonnegative for nonnegative initial conditions.

To prove the nonnegativity of solutions of (3.1)-(3.2), let $\left\{\psi_1^{(n)}(s)\right\}, s \in [-\tau, 0], i = 1,2; (n = 1,2, \ldots)$ be a sequence of strictly positive continuous functions such that

$$\lim_{n\to\infty} \psi_1^{(n)}(s), = \psi_1(s), \qquad s \in [-\tau, 0], \qquad i = 1,2 \tag{3.3}$$

Let $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ be the solution of (3.1) corresponding to the initial condition

$$\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \left\{\psi_1^{(n)}(s), \psi_2^{(n)}(s)\right\}, \quad s \in [-\tau, 0],. \tag{3.4}$$

Consider the solution $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ for $t \in [0, \tau]$ (assuming it exists on the whole of $[0, \tau]$). Suppose $x_1^{(n)}$ does not remain nonnegative for all $[0, \tau^*]$; then there exists a $t^* = [0, \tau^*](\tau^* =$ the positive minimum of $\tau_{ij}$ $(i, j = 1, 2))$ for which

$$x_1^{(n)}(t^*) = 0 \ and \ x_1^{(n)}(t)(t^* - \tau_{11})) > 0.$$

It will follow from (3.1) and the positivity of the initial condition that

$$\frac{dx_1^{(n)}(t^*)}{dt} = b_2\left(x_1^{(n)}(t^* - \tau_{11})\right) > 0 \tag{3.5}$$

Where we have used the properties of $b_i$ and $m_i$ from (2.2)-(2.6); it is found that (3.5) contradicts the definition of $t^*$. Thus $x_1^{(n)}(t) > 0$ on $[0, \tau^*]$; similarly $x_2^{(n)}(t) > 0$ $on$ $[0, \tau^*]$. We can repeat this procedure for intervals of the form $[\tau^*, 2\tau^*], [2\tau^*, 3\tau^*]$ $etc$. Thus it will follow that so long as $(x_1^{(n)}(t), x_2^{(n)}(t))$ exists, we have

$$x_1^{(n)}(t) \geq 0 \, x_2^{(n)}(t) \geq 0.$$

If we now consider the limits as $n \to \infty$, we get

$$\lim_{n\to\infty}\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \{x_1(t), x_2(t)\} \tag{3.2}$$

With $x_1(t) \geq 0, , x_2(t) \geq 0$ where $(x_1(t), x_2(t))$ is the solution of (3.1)-(3.2) and this is a consequence of the continuous dependence of solutions of (3.1) on initial conditions (Hale [17], page 41).

A second question for (3.1) with (2.2)-(2.6) is whether solutions of (3.1) exist for all $t > 0$? suppose now that solutions of (3.1) do not exist for all $t$ ; then there exists a $t_1 > 0$ such that at least for one of either of $x_1$ *or* $x_2$ we have

$$\lim_{t \to t_1 -} x_i\ (t) = \infty. \tag{3.7}$$

To be specific let us suppose that $\lim_{t \to t_1 -} x_i\ (t) = \infty$; let $t_2$ be the first instant of time for which $x_1(t_2) = \delta_1, t_2 < t_1$. It will follow from (3.1) that

$$\frac{dx_1(t_2)}{dt} = b_1\big(x_1\ (t_2 - \tau_{11})\big) - m_1\big(x_1(t_2), x_2\ (t_2 - \tau_{12})\big)$$

$$< b_1\big(x_1\ (t_2)\big) - m_1\big(x_1(t_2), x_2\ (t_2 - \tau_{12})\big)$$

$$< b_1(\delta_1) - m_1\big(\delta_1, x\ (t_2 - \tau_{12})\big) < 0 \tag{3.8}$$

Which contradicts the definition of $t_2$. Thus solutions of (3.1), (2.2) – (2.6) exist for all $t$ .

We are now ready to consider the effects of time delays in reproduction and interspecific interactions on the coexistence of the two competing species. The steady state $(\alpha, \beta)$ of (2.1) is again a steady state of (3.1). To examine the (local) asymptotic stability of $(\alpha, \beta)$ in (3.1) we let

$$x_1(t) = \alpha + Y_1(t), \quad x_2(t) = \beta + Y_2(t)$$

In (3.1) and derive that the linear variational system governing the perturbations $Y_1, Y_2$ is given by

$$\frac{dY_1\ (t)}{dt} = -m_{11}Y_1(t) + b_{11}Y_1\ (t - \tau_{11}) - m_{12}Y_2\ (t - \tau_{12}),$$

$$\frac{dY_2\ (t)}{dt} = -m_{22}Y_2(t) + b_{21}Y_1\ (t - \tau_{21}) - m_{22}Y_2\ (t - \tau_{22}),$$

Where

$$m_{ij} = \frac{\partial m_1}{\partial x_j}, \quad b_{ii} = \frac{\partial b_i}{\partial x_i}\ at\ (\alpha, \beta). \tag{3.10}$$

It is well known (Bellman and Cooke [3], page 336) that the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) if the trivial solution of the variational system (3.9) is asymptotically stable. We formulate our first result in the following form whose proof is given in the appendix.

*THEOREM 3.1. Assume that in the two systemic types competition model system (3.1) with delayed density dependent birth and death rates $b_i$ and $m_i$ $(i = 1, 2)$ satisfy the conditions (2.2)-(2.6). Suppose that the conditions (2.10) hold. Then the nontrivial steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) whatever the delays be $\tau_{ij}$ $(i = 1, 2)$.*

Let $y_1(t) and\ y_2(t)$ respectively denote the population sizes (or biomasses) of a prey species and a predator species which feeds exclusively on the single prey species in a temporally constant environment. In the format of (1.5) the following system will describe the rates of change of the two population sizes:

$$\frac{dy_1(t)}{dt} = f_1(y_1(t)) - g_1(y_1(t), y_2(t)),$$
$$\frac{dy_2(t)}{dt} = f_2(y_1(t)) - y_2(t) - g_2(y_2(t)), \quad (4.1)$$

Where $f_i\ and\ g_i\ i = 1,2)$ respectively denote the birth and death rates of the two species; in order to make (4.1) denote the dynamics of a prey – predator species we will assume that $f_i\ and\ g_i$ satisfy the following conditions (see for instance Rescigno and Richardson [26]); $f_i\ and\ g_i$ are continuous with continuous partial derivatives for all $y_i \geq 0\ (i = 1, 2); also\ we\ require$

$$\frac{\partial f_1}{\partial y_1} > 0; \quad \frac{\partial g_1}{\partial y_1} \geq 0; \quad \frac{\partial g_1}{\partial y_2} \geq 0,$$
$$\frac{\partial f_2}{\partial y_1} \geq 0; \quad \frac{\partial g_2}{\partial y_2} \geq 0; \quad \frac{\partial g_2}{\partial y_2} \geq 0, \quad \text{Where } y_1 \geq 0, y_1 \geq 0, \quad (4.2)$$

(i)

$$f_1(0) = 0, \quad g_1(0, y_2 \equiv 0,$$
$$f_2(0, y_2) = 0 = f_2(y_1, 0); \quad g_2(0) = 0, \quad (4.3)$$

(ii) For some constant $y_1^* > 0$ we have

$$f_1(y_1^*) - g_1(y_1^* 0) = 0, \quad (4.4)$$

(iii) There exists a point $(\alpha^*\beta^*), \alpha^* > 0, \beta^* > 0$ such that

$$f_1(\alpha^*) - g_1(\alpha^*, \beta^*) = 0,$$
$$f_2(\alpha^*, \beta^*) - g_2(\beta^*) = 0, \quad (4.5)$$

(iv) There exist positive constants $n_1, n_2$ such that

$$f_1(\eta_1) - g_1(\eta_1, y_2) < 0, y_2 \geq 0,$$
$$f_2(\eta_1, \eta_2) - g_2(\eta_2) < 0 \quad (4.6)$$

The conditions (4.2) mean that the system (4.1) is of the oppugnacy systems type; (4.3) will imply that (0,0) is a trivial steady state of (4.1); (4.4) means that in the absence of the predator, the prey species has a nontrivial (positive) steady state;

(4.5) implies that (4.1) has a nontrivial steady state in the interior of the positive quadrant of the state species of (4.1); (4.6) can be shown (see below) to imply that no credit type and debit type system types can grow unboundedly.

An example of (4.1)-(4.6) is the familiar Volterra-Lotka prey-predator system

$$\frac{dy_l(t)}{dt} = r_l y_l(t) - [a_{ll} y_l^2(t) + a_{12} y_l(t) y_2(t)],$$
$$\frac{dy_2(t)}{dt} = a_{12} y_1(t) y_2(t) - [r_2 y_2(t) + a_{22} y_2^2(t)],$$

Where $r_{l,}\ a_{ij}\ (i,j=1,2)$are positive costants such that

$$a_{21}r_1 - r_2a_{11} > 0.$$

As in the previous case of competition we have to verify that when $y_1(0) > 0, y_2(0) > 0$ the corresponding solution of (4.1) will be nonnegative for $t \geq 0$ for which the solution exist. Since the co-ordinate axes $y_1 = 0, y_2 = 0$ are invariant sets for (401) by virtue of (4.3), it will follow that the solution of (4.1) starting from the interior of the positive quadrant cannot reach the outside of that quadrant. This is a consequence of the invariance of the co-ordinate axes for (4.1)

To prove that solutions of (4.1)-(4.6) exist for all t $\geq 0$ let us suppose that $y_l 0\ and\ y_2(0) satisfy\ 0 < y_l\ (0)$ satisfy $0< y_l(0) < \eta_l, 0 < y_2(0) < \eta_2$ ( on the other hand if either one of $y_l(0) \geq \eta_l and\ y_2(0) \geq \eta_l$ or both hold, it will follow that the corresponding $y_l(t)$ will decrese by virtue of (4.6) at least for small $t \geq 0,$ ). Suppose $y_l(t)$ does not exist for all $t \geq 0$; then for some $t_l < \infty\ lim_{t\to t_l-} y_l(t) = \infty$and let $t_2 > 0$ be the first instant for which $y_1(t_2) = \eta_1$; then we have from (4.1) and (4.6) that

$$\frac{dy_1(t_2)}{dt} = f_1(\eta_1) - g_1(\eta_{1,}\ y_2) < 0\ (assumimg\ y_1(0) < \eta_1)$$

Which contradicts the definition of $t_2$ showing that $y_1(t)$ exists for all $t \geq 0$; if $y_2(t)$ is not defined for all $t \geq 0$, then for some $t_3 \geq \infty$ we will have $lim_{t-t_3-} y_2(t) = \infty$;now let $t_4$ be the first instant for which $y_2(t_4) = \eta_2$; then from the properties of $f_2$ we have

$$\frac{dy_2(t_4)}{dt} = f_2(y_1(t_4), \eta_2) - g_2(\eta_2) \leq f_2(\eta_1, \eta_2) - g_2(\eta_2)$$

$$< 0 (by\ (4.6) assuming\ 0 < y_1(0) < \eta_1, 0 < y_2(0) < \eta_2)$$

Which again contradicts the definition of $t_4$. Thus $y_2(t)$ is defined for all $t \geq 0$.

We will now derive a set of sufficient condition for the (local) asymptotic stability of the steady state $(\alpha^*, \beta^*)$ of (4.1). in the next section we will show that under the same conditions and in the presence of various delays , the steady state $(\alpha^*, \beta^*)$remains (locally) asymptotically stable.

To examine the asymptotic stability of the nontrivial steady state $(\alpha^*, \beta^*)$ in (4.1) we let

$$y_1(t) = \alpha^* + U_1(t), \qquad y_2(t) = \beta^* + U_2(t)$$

In (4.1) and derive the liner variational system in the perturbations $U_1, U_2$ as follows:

$$\left.\begin{aligned} \frac{dU_1(t)}{dt} &= \frac{\partial f_1}{\partial y_1}U_1(t) - \frac{\partial g_1}{\partial y_1}U_1(t) - \frac{\partial g_1}{\partial y_2}U_2(t), \\ \frac{dU_2(t)}{dt} &= \frac{\partial f_2}{\partial y_1}U_1(t) - \frac{\partial f_2}{\partial y_2}U_2(t) - \frac{\partial g_2}{\partial y_2}U_2(t), \end{aligned}\right\} \quad (4.7)$$

Where all the partial derivatives in (4.7) are evaluated at $(\alpha^*, \beta^*)$. The steady state $(\alpha^*, \beta^*)$ of (4.1) is (locally) asymptotically stable if and only if the trivial solution of (4.7) is asymptotically stable which will be the case if and only if the roots of the characteristic equation

$$\text{Det} \begin{pmatrix} \lambda - (f_{11} - g_{11}) & g_{12} \\ -f_{21} & \lambda - (f_{22} - g_{22}) \end{pmatrix} = 0,$$

Where $f_{ij} = \frac{\partial f_1}{\partial y_l}$, $g_{ij} = \frac{\partial g_1}{\partial y_j}$, at $(\alpha^*, \beta^*)$ (4.8)

$(i,j = 1, 2)$

Have negative real parts. If $\lambda_1, \lambda_2$ are the riits of (4.8) then

$$\lambda_1 + \lambda_2 = (f_{11} - g_{11}) + (f_{22} - g_{22}), \quad (4.9)$$

$$\lambda_1 \lambda_2 = (f_{11} - g_{11})(f_{22} - g_{22}) + f_{21} g_{12}.$$

The proof of the following result is an immediate consequence of (4.7)-(4.9)

THEOREM 4.1 *In the CREDIT-DEBIT system (4.1)-(4.6) assumes that the partial derivatives satisfy the following additional conditions:*

$$\frac{\partial g_1}{\partial y_1} > \frac{\partial f_1}{\partial y_1} + \frac{\partial f_2}{\partial y_1}, \quad \frac{\partial g_2}{\partial y_2} > \frac{\partial f_2}{\partial y_2} + \frac{\partial g_1}{\partial y_2}, \quad \text{Evaluated at } (\alpha^*, \beta^*), \quad (4.10)$$

*Then the steady state $(\alpha^*, \beta^*)$, is (locally) asymptotically stable for the prey- predator system.*

## 5. Delays in reproduction, predation and gestation

A number of authors have proposed and analysed several models of prey- predator systems with time delays. Most of the models have been derived from the so called models of the Kolmogorov-type f (1.2) with time delays incorporated in the average growth rates of the prey and predator species. Volterra [32] proposed the following model for a prey-predator system. The equations are used for the formulation of Credit-debit system of atoms, accounts, economic variable, cosmological variable, astrophysical; variables, and qubit systems for which the model is being applied.

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \delta_1 N_1(t) - \int_{\infty}^{t} F_1(t-s) N_2(s) ds],$$

(5.1)

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 + \delta_2 N_1(t) - \int_{\infty}^{t} F_2(t-s)N_1(s)ds],$$

Where $r_1, r_2, \delta_1, \delta_2$ *are positive constant and* $F_1, F_2$ is nonnegative continuous delay kernel defined and integrable on $[0, \infty]$. (5.1) has no predator self-regulating negative feedback. Brelot [6] considered a modified format of (5.1) in the form

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \lambda_1 N_1(t) - \delta_1 N_2(t) - \int_{-\infty}^{t} F_1(t-s)N_2(s)ds],$$

(5.2)

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 - \lambda_2 N_2(t) - \delta_2 N_1(t) - \int_{-\infty}^{t} F_2(t-s)N_1(s)ds],$$

Which has term representing intraspecific competition among both predator and prey species. Under suitable condition on the various parameters in (5.2) one can show that (5.2) has a (locally) asymptotically stable steady state. A number of intergrodifferential equation models of the types (5.1) and (5.2) have been investigated by Cushing [12]

Starting from Hutchinson's (1948)[19] delay logistic equation, May[23] has proposed the following system

$$\frac{dN_1(t)}{dt} = rN_1(t)\left[1 - \frac{N_1(t-\tau)}{K}\right] - \alpha N_1(t)N_2(t),$$

$$\frac{dN_2(t)}{dt} = -bN_2(t) + \beta N_2(t)N_1(t), \qquad (5.3)$$

Where r, $\tau$, K, $\alpha, \beta, b$ are positive constants; (5.3) contains a single discrete delay; one can modify (5.3) and incorporate a continuously distributed delay in it so that (5.3) can be written as (May [23])

$$\frac{dN_1(t)}{dt} = rN_1(t)[1 - \int_{-\infty}^{t} F(t-s)N_1(s)ds] - \alpha N_1(t)N_2(t),$$

$$\frac{dN_2(t)}{dt} = -bN_2(t) + \beta N_1(t)N_2(t). \qquad (5.4)$$

The first model of a prey-predator type which departs from the Kolmogorov-type formulation is due to Wangersky and Cunningham [33] who considered the model system

$$\frac{dN_1(t)}{dt} = a_1N_1(t) - b_1N_1^2(t) - c_1N_1(t) - N_2(t),$$

$$\frac{dN_2(t)}{dt} = a_2N_2(t) + c_2N_1(t-\tau)N_2(t-\tau), \quad (5.5)$$

Where $a_1, a_2, b_1, c_1, c_2, \tau$ are positive constants; (5.5) means that a duration of $\tau$ time units elapses when an individual prey is killed and the moment when the corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) Recently Cushing [13] has considered a model of the form

$$\frac{dN_1(t)}{dt} = rN_1(t)\left[1 - \frac{N_1(t)}{k} - \alpha N_2(t)\right],$$

$$\frac{dN_2(t)}{dt} = -d_2^*(t) + bN_1(t)\int_0^\infty \delta(a)N_2(t-a)e^{-d*a}\,da, \quad (5.6)$$

Where $r, k, \alpha, d^*, b$ are positive constants and $\delta$ is related to the age dependent fecundity of the predator species.

With the models (5.1)-(5.6) in the background let us now consider a time delayed model similar to (4.1) in the following form:

$$\frac{dN_1(t)}{dt} = f_1\big(N_1(t-\tau_{11})\big) - g_1\big(N_1(t), N_2(t-\tau_{12})\big),$$

$$\frac{dN_2(t)}{dt} = f_2\big(N_1(t-\tau_{21})\big), N_2(t-\tau_{12}) - g_1(N_2(t)), \quad (5.7)$$

Where $\tau_{ij}$ $(i,j=1,2)$ are nonnegative constants with $\tau = \max \tau_{ij}, i,j = 1,2$; the birth and death rate functions $f_1, f_2, g_1, g_2$ satisfy the same conditions as in (4.1)-(4.6). Along with (5.7) we specify the initial conditions as follows:

$N_1(s) = \phi_1(s) > 0,$

$s \in [-\tau, 0].$

$$N_1(s) = \phi_1(s) > 0,$$

We3 will assume that in addition to (4.2) we have for (5.7) that

$$\frac{\partial f_2}{\partial N_1} > 0, \frac{\partial f_2}{\partial N_2} > 0, f_2(N_2, N_2) > 0 \; whenever \; N_1 > 0, \; N_2 > 0. \; (5.9)$$

It is not difficult to show by the methods previously used for the competition system that whenever $N_1(s) > 0$ and $N_2(s) > 0$ for $s \in [-\tau, 0]$, solution of (5.7) exist for all $t \geq 0$ and remain nonnegative for all $t \geq 0$. also $(\alpha^*, \beta^*)$ of (4.5) is a steady state of (5.7).

We will discuss the form and meaning of the delays in (5.7) in the next section. The effects of the various delays in (5.7) on the (local) asymptotic stability of $(\alpha^*, \beta^*)$ can be derived from the following theorem whose proof is quite similar to that of Theorem 3.1 and hence the details of the proof are omitted.

THEOREM 5.1 *In the time delayed prey-predator system(* corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) *) (5.7) let the conditions (4.2)-(4.6), (5.8) and (5.9) hold for the birth and death rates; furthermore suppose the condition (4.10) hold. Then for all nonnegative delays in (5.7), its nontrivial steady state $(\alpha^*, \beta^*)$ is (locally) asymptotically stable.*

## 6. Estimation of decay rates

We will first derive a convention sufficient condition for the asymptotic stability of the trivial solution of a linear system of the form

$$\left.\begin{aligned} \frac{dY_1(t)}{dt} &= -m_{11} Y_1(t) + b_{11} Y_1(t-\tau_{11}) - m_{12} Y_2(t-\tau_{12}), \\ \frac{dY_2(t)}{dt} &= -m_{22} Y_1(t) + m_{21} Y_1(t-\tau_{21}) - b_{22} Y_2(t-\tau_{22}), \end{aligned}\right\} \; t > 0, \quad (6.1)$$

Where $m_{1J}, b_{i1}$ $(i, j = 1,2)$ are positive constants while $\tau_{1J}(i, j = 1,2)$ are non negative with $\tau = \max(\tau_{1J})(i, j = 1,2$.as is shown in the appendix, the following system dominates (6.1*)*

$$\left.\begin{aligned} \frac{dZ_1(t)}{dt} &= -m_{11} Z_1(t) + b_{11} Z_1(t-\tau_{11}) - m_{12} Z_2(t-\tau_{12}), \\ \frac{dZ_2(t)}{dt} &= -m_{22} Z_1(t) + m_{21} Z_1(t-\tau_{21}) - b_{22} Z_2(t-\tau_{22}), \end{aligned}\right\} \; t > 0, \quad (6.2)$$

$Z_i(s) = |Y_i(s)|, \qquad i = 1,2 \;\; for \; s \leq 0 \;\; for \; which \; Y_i(s) \; is \; defined.$

Thus the asymptotic stability of the trivial solution of (6.2) implies that of (6.1). it is also shown in the appendix that the solutions of (6.2) are nonnegative for all t $\geq 0$. we will use the following vector-matrix notation:

$$Z = (Z_1, Z_2)^{\tau}, \qquad ||Z|| = |Z_1(t)| + |Z_2(t)| = Z_1(t) + Z_2(t)$$

$$B = \begin{bmatrix} b_{11} & m_{12} \\ m_{21} & b_{12} \end{bmatrix}, \quad ||B|| = max \begin{matrix} b_{11} + & m_{21} \\ b_{21} + & m_{12} \end{matrix}, \tag{6.3}$$

$$A = \begin{bmatrix} -m_{11} & 0 \\ 0 & -m_{22} \end{bmatrix}.$$

The logarithmic norm or the measure of any square matrix B, denoted by $\mu(B)$is defined as follows:

$$\mu(B) = \lim_{\theta \to 0+} \frac{/\!/I + \theta B/\!/ - 1}{\theta}$$

($I$ being an identity matrix of the same size as $B$).The matrix measure compatible with the norms in (6.3) is given by

$\mu(B) = \max_J [b_{jj} + \sum_{\substack{i=1 \\ i \neq J}}^{2} | b_{jJ} | ]$ $B=(b_{1J})$.

We refer to Coppel [11]or Vidyasager [31]for a number of properties of matrix measure We have shown in the appendix that the derivative of $H(Z_1,Z_2)(t)$(seeA.10)satisfies the relation

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq [-m_{11}+b_{11}+m_{21}]Z_1(t)+[-m_{22}+b_{22}+m_{12}]Z_2(t)$, and hence we have

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq \{\mu(A) + /\!/B/\!/\}\{ Z_1(t)+ Z_2(t)\}$

If we assume for the matrix norm $\|B\|$ as in (6.3)that

$$\mu(A) + /\!/B/\!/ < 0 \; or \; -\mu(A) > /\!/B/\!/$$

Then from the results of the appendix (Lemmas 5 and 6) it will follow that $Z_1(t)+Z_2(t) \to 0$ as $t \to \infty$ and thus we have the following:

Lemma 6.1 Assume that the Coefficients of the linear system (6.1)satisfy the condition (6.7). Then the trivial solution of(6.1)is asymptotically stable.

When (6.7)holds an estimate for the decay rate of solutions of (6.1)can be obtained as follows A necessary and sufficient condition for the asymptotic stability of the trivial solution is that all the roots of the characteristic equation

Det $G(\lambda,T_{iJ})=0$

Have negative real parts where

$G(\lambda,T_{1J})=[\begin{matrix} \frac{\lambda+m_{11}}{0} & \frac{0}{\lambda+22} \end{matrix}] - b_{11}^{e-\lambda T_{11}}$ $\begin{matrix} -m_{12}^{e-\lambda T_{12}} \\ -m_{21}^{e-\lambda T_{21}} \quad b_{22}^{e-\lambda T_{22}} \end{matrix}$

$=\lambda I - A - B(\lambda, T_{ij})$(say),

A positive number $\mathfrak{a}$(if it exists)is said to be a decay rate of solution of (6.1) if and only if every root of (6.8)satisfies the condition

$\mathrm{Re}(\lambda)+\mathfrak{a} < 0.$

When (6.11)holds the solution of (6.1)will be of the order of $e^{-at}$ and hence the supremum of all such $\mathfrak{a} > 0$ $a$ satisfying (6.11)will be a desirable parameter to estimate However such a supremum of all such (6.11)may not be possible to determine since it has to be done directly from (6.8)-(6,10)(see also Section 3).We will obtain an estimate of $\mathfrak{a}$ (not necessarily the supremum)I (6.11)by the following procedure.

Replace $\lambda$ byz-$\mathfrak{a}$ in (6.8)and define

$G(z-\mathfrak{a}, T_{IJ}) = M(z) = (z-a)I - A - B(z-\mathfrak{a}, T_{IJ}).$

By Lemma 6.1 a Sufficient condition for all roots of M(z)=0 to have negative real parts equivalently to satisfy

$\mathrm{Re}(\lambda) < -a$

is the following :

$-\mu(aI + A) > \|B_1(a, T_{ij})\|,$

Where

$$B_1(a, T_{ij}) = \begin{pmatrix} b_{11}^{e-\lambda T_{11}} & -m_{12}^{e-\lambda T_{12}} \\ -m_{21}^{e-\lambda T_{21}} & b_{22}^{e-\lambda T_{22}} \end{pmatrix}$$

Since $\mu(\mathfrak{a} I + A) = \mathfrak{a} + \mu(A)$ we have from (6.13) that

$\mathfrak{a} + \mu(A) + \|B_I(\mathfrak{a}, T_{IJ})\| < 0$

Is a sufficient condition for Re $(\lambda) < -a$If we denote the left side of (6.15)by $\mathfrak{a} f(\mathfrak{a}, T_{Ij})$ we note that

$f(0, T_{Ij}) = \|B\| + \mu(A) < 0 \ \ if \ (6.7) holds;$

Al so $\partial f / \partial \mathfrak{a} > 0, f$ is continuous and monotonic in $(\mathfrak{a})$;$\to \infty as\ \mathfrak{a} \to \infty$; thus by the continuity of $f$ on $\mathfrak{a} \geq 0$, there exists an interval of the from $[0, ^{*})$such that

$f(0, T_{Ij}) < 0 \ \ for\ \mathfrak{a} \in [0, \mathfrak{a}^{*})$

And hence

$f(\mathfrak{a}, T_{Ij}) = 0$

Thus an upper bound $\mathfrak{a}^*$ of $\mathfrak{a}$ for which we have $f(\mathfrak{a}, T_{Ij}) < 0$ $\mathfrak{a} \in [0, \mathfrak{a}^*)$ can be obtained as the unique (by the monotonicity of $f$ on $\mathfrak{a}$) positive root of$\max \left\{ b_{11}^{e-\lambda T}{}_{11} - m_{12}^{e-\lambda T}{}_{21} + \mathfrak{a}\text{-}\min \left\{ m_{11} \right.\right.$

$\left.\left. b_{22}^{e-\lambda T}{}_{22} - m_{21}^{e-\lambda T}{}_{12} \right\} \right\} m_{22}$

Where the maximum in (6.18) is taken over $b_{11}, b_{22}, m_{12}, m_{21}$ and $T_{ij}$ $(i, = j\ 1,2)$. A numerical procedure can be devised to solve equations of the form (6.18) in concrete cases. If one is interested in a smaller bound on the decay rate it can be obtained as the unique root of

$$\| B \| e^{ar} + a + \mu(A) = 0, \tau = \max_{i,J} T_{ij}. \qquad (6.19)$$

It will follow from (6.19) than an increase in T

Will reduce the value of the root of (6.19). We remark that since our estimate $a^*$ of dacay rate is based on a sufficient condition of the form (6.7), we can conclude that the solotions of (6.1) will converge at least as fast as $e^{-a*t}$ where $a^*$ is the unique root of either of (6.18) or (6.19).

## Implications

Authors have formulated a class of two systems model systems staring from the elementary principal that the rate of change of a population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) equals the creation rate minus the destruction rate. This is a departure from the conventional model building in population ecology where one usually starts from the Malthusian formulation, now known as the Kolmogorov formulation in the nonlinear context. In our analysis of the models proposed we have obtained sufficient conditions for the global existence of nonnegative solutions of the model systems (this is trivial in the Kolmogorov formulation) and the other desirable properties such as the existence of suitable steady states. The two classes of models proposed for competition and debt contain intraspecific self – regulation biologists that each systems is under the action of a self-regulating negative feedback force, opinions and arguments differ regarding the mechanism by which such an effect is realized; we refer for more details on this aspect to a collection of articles compiled by Tamarin [30].

Gopalasamy's main result is the following : if a self –regulating negative feedback mechanism (intraspecific oppugn ant) can act without time delay and if such an effect can keep an ecosystem ( a community of one or more species of populations) asymptotically stable ( at least locally about a non trivial steady state, then time delays of any magnitude in any other interaction involving positive or negative feedbacks cannot destabilize such a stable ecosystem ( in a mathematical sense); that is such delays cannot cause a hopf-type bifurcation leading to persistent and undamped fluctuations. The delay-induced bifurcations to oscillations in models have been investigated by many authors. (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Qualitatively result is not new to the literature the effects of delays on otherwise stable and unstable systems have been considered by Cushing[12]where (pages 35-39)he has shown that an otherwise unstable system can be made

less unstable by delays, Such an effect of delays on unstable systems appear to have been first proposed by Beddinton and May [2]who have also shows that delays can make stable systems less stable leading to a conclusion "time delays are not necessarily destabilizing" It is also shown by Cushing([12], pages 75-103)that while time delays cannot stabilize(in the mathematical sense)an otherwise unstable two species system in competition or mutualism time delays in there species systems can stabilize otherwise unstable systems (even in a mathematical sense)Cushing and Saleem [14]have formulated a model where delays in the positive feedback can have a stabilizing effect: Hale[18]discusses briefly the stabilizing nature of delays and given an example to show that an increase in delays can make an unstable system stable Thus for a better understanding of the effects of delays on otherwise stable systems one has to estimate the rate of decay of perturbations (or the characteristic return time )in terms of the parameters of the system especially when the delays do not cause a mathematical instability This aspect has been investigated by a procedure slightly different from that of the characteristic equation we have proposed (section6) method for the estimation of decay rate of perturbations when there are multiple delays Since many species systems with multiple delays cannot be analysed (due to complexity)by the characteristic equation our methods using comparison lead to easily verifiable sufficient conditions The sufficient conditions obtained are in the form of a diagonal dominance type; these conditions can be weakened to quasi-diagonal dominance type as in Ladde [20], Siljak[20]Anderson [1], Lewis and Anderson [21,22](Ladde's result [20] is delay dependent due to the fact that delay appears in self –regulation also).Verification of quasi –diagonal dominance conditions in arbitrary matrices (except for Metzler matrices )is an impossible task and quasi –diagonal dominance is hard to interpret.

If follows from our discussion that those confederation systems which display delay induced instability leading to permanent and undamped oscillations should have one or more of the following aspects:

(1)time delays are in intraspecific interaction terms (and possibly in other terms also): (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

(2) Intraspecific self-regulation terms are absent in density dependent from in one or more populations of the system;

(3) Intraspecific negative feedback effect are not strong enough to maintain the stability of the relevant steady state of the nondelayed system.

We can now infer how the delay induced stability leads to fluctuations; a primary mechanism of such fluctuations can be described at least theoretically as follows : in its attraction of a system to a steady state, intraspecific self-regulation fails to be effective and the system overhoots the steady state and at such higher densities intraspecific competition becomes operationally effective leading the system to understood the steady state and this process is repeated again and again resulting in undamped oscillations about the steady state. It follows from our results that if intraspecific competition is sufficiently strong, the system cannot become oscillatory even when there are time delays, so long as they do not interfere with the strength of the intraspecific self – regulation.

We are thus led to the question of existence and importance of intraspecific competition in population systems. In the introduction to an anthology of articles by several authors, Tamarin [ 30] begins with the statement "all animal populations have the reproductive capacity to increase without limit, yet none di" and in his concluding remarks writes " this leaves us with only one great concept in the area of confederation do main regulation: that all credit-debit are regulated" More specifically this is applicable to all the confederation of (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions,

financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

In his discussion on the effect of competition on the densities of populations, Nicholson ([ 25], $page$ [ 147] ) has reached the conclution "Intraspecific competition automatically regulates the severity of its action to the requirements of each denizen , provided the inherent resistance of the environment is sufficiently low to permit the species to exist." In order to look for a mechanism by which intraspecific competition plays its role of decreasing growth rate as confederation density increases, Calhoun [ 7] set up an experiment with Norway rates (Rattus norvegicus) and found that increased density leads to increased hostility which disturbs the psychology and physiology causing decreased reproductive output. Many other mechanisms by which intraspecific compition self-regulate populations have been put forward by a number of authors (see Tamarin [30]). This is true of (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

If we accept the view that intraspecific competition is really predominant in population regulation, then we can understand how and why time delays destabilize stable system leading to oscillations; time delays render intraspecific competitive regulation territorially (in case of particles we take the field in which it operates) weak so that self –regulatory forces become temporarily insufficient. If we now agree that all population models should incorporate self-regulating intraspecific competition terms in surfeit strength along with other tropic intraspecific competition terms in sufficient strength along with other atrophic interactions then it is not fair to say " time delays destabilize" and we note from our analysis that it is usually a failure in self- regulation which can cause instability rather than delays by themselves.

We will give one example to illustrate our point. Consider a credit-debit system in which the debits is self –regulating by a logistic load as follows: (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

$$\frac{dx}{dt} = r\left(x - \frac{x^2}{K}\right) - axy = b_I(x) - m_I(x,y) \quad (say),$$

$$\frac{dy}{dt} = \beta xy - \delta y = b_2(x,y) - m_2(y) - \quad (say),$$

Where $K, a, \beta, \delta$ are positive constants such that the steady state $x^*\ y^*$ exists where

$$x^* = \frac{\delta}{\beta}; \quad y^* = \frac{r}{a}\left(-\frac{\delta}{K\beta}\right) > 0.$$

In this example the predator self –regulation $(\partial m_2/\partial y)$ is not sufficiently strong in the sense

$$x^* = \left(\frac{\partial m_2}{\partial y}\right)_{(x^* y^*)} \succ \left(\frac{\partial b_2}{\partial y} + \frac{\partial m_1}{\partial y}\right)_{(x^* y^*)}$$

If in such a model there is a time delay in the credit creation rate, such a delay can lead to instability of $(x^*\ y^*)$ of setting up undamped oscillations. If however we modify the growth rate by the inclusion of a density dependent self-regulation so that credit

$$\frac{dy}{dt} = \beta\, xy - \delta y - \eta y^2$$

For some positive constant η then it is possible, depending on η, for such a modified system to remain stable without oscillations even if there is a time delay in the credit creations rate. We have another completely different justification to have $\eta y^2$ in the credit dynamics: for instance if the debit is removed from this cannot happen. While this example serves to illustrate an aspect of our result, our advocacy of the significance is not based on this example.

Gopalasamy has thus derived conclusion that "time delays need not destabilize" in models with several discrete delays. However we have provided in the appendix the necessary tools (Lemmas 3,4,5) to prove all our results when the relevant delays are all (or some) continuously distributed. Some authors (Caswell [9], May [23], Cushing [12] appear to regard continuously distributed dalays as financially/cosmologically credit wise obtained when he subjected algae to a variable nitrate environment and obtained a better fit for his data with a continuous delay. The questions of whether a discrete or continuous delay is more realistic may well depend on what process is being modeled. For instance an increased supply of some resource at some point in time may trigger reproduction at a particular moment in time later irrespective of the conditions of the resource during the intervening period. If the duration of increased resource supply necessary to elicit such a response is short compared to the intervening period before the reproduction occurs, then a discrete time delay may be appropriate. In a large number of cases of animal. Note fungus algae are prey. Here it shall be an account which loses the energy, amount, or matter for that matter. Same result holds good for (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Populations of particles, atoms, molecules, propositions, systems types chosen the past dietary and nutritional history of an animal over a long period plays an influential role in determining the current bahaviour of the dietary and nutritional history of an animal over a long period plays an influential role in determining the current behavior of the animal and in such cases continuous time delays will be appropriate.

In conclusion the following comments are appropriate. The time delays since the current creation rate is usually due to a contribution of those mature animals living a few units of time earlier depending on the type of animal or organism; the time delays appearing in interspecific interactions can be due to the fact that in order to successfully interact and compete interferentially a developmental maturity may be necessary. In the case of predation, delays in the birth rate of prey can be due to delayed maturation while that in the birth rate of a predator is due to gestation delay with delayed fecundity; the delays in the predation rate are again are again as before due to the maturity needed for a predator to successfully organize and succeed in an attack.

In our analysis we have assumed throughout a temporally uniform environment neglecting dispersal of populations and their age structure. We expect our results to hold under suitable conditions both in temporally periodic environments with the steady state of a constant environment being replaced by a periodic solution

(nonstationary steady state) and when the effects of dispersion and age structures are considered. It is worthwhile to modify our purely ecological, cosmological, financial, propositional, problematic systems models so as to provide them with a genetic foundation and examine the effects of time delays since forces of selection act with some time delay. Such a generalization will definitely lead to a better understanding of "how confederations regulate themselves".

## Appendix

First we derive a comparison result.

Lemma 1. Consider the delay differential system

$$\left.\begin{aligned}\frac{dY_1(t)}{dt} &= -m_{11}\ Y_1(t) + b_{11}Y_1(t-\tau_{11}) - m_{12}\ Y_2(t-\tau_{12}),\\ \frac{dY_2(t)}{dt} &= -m_{22}Y_2(t) - m_{21}Y_1(t-\tau_{21}) + b_{22}\ Y_2(t-\tau_{22}),\end{aligned}\right\} (A.1),$$

Where the $m_{ij}$, $b_{ij}$ $(i.j = 1,2)$ are positive while $\tau_{ij}$ $(i.j = 1,2)$ nonnegative with $\tau = \max \tau_{ij}$, $i,j = 1,2$. Then

$$\left.\begin{aligned}\frac{d|Y_1(t)|}{dt} &\le -m_{11}\,|Y_1(t)| + b_{11}|Y_1(t-\tau_{11})| + m_{12}\ Y_2(t-\tau_{12})|,\\ \frac{d|Y_1(t)|}{dt} &\le -m_{22}\,|Y_2(t)| + m_{21}|Y_1(t-\tau_{21})| + b_{22}\ |Y_2(t-\tau_{21})|,\end{aligned}\right\} (A.2)$$

Where

$$\frac{d}{dt}\,|\,Y_1(t)\,| = \left\{\begin{aligned}&\frac{dY_1(t)}{dt}\ if\ \ Y_1(t) \ge 0,\\ &\frac{dY_1(t)}{dt}\ if\ \ Y_1(t) < 0,\end{aligned}\right\}\ i = 1,2. \qquad (A.3)$$

The proof is by direct verification.

LEMMA 2. Consider now (A.2) and the following system

$$\left.\begin{aligned}\frac{dZ_1(t)}{dt} &= -m_{11}\ Z_1(t) + b_{11}Z_1(t-\tau_{11}) + m_{12}\ Z_2(t-\tau_{12}),\\ \frac{dZ_2(t)}{dt} &= -m_{22}Z_2(t) + m_{21}Z_1(t-\tau_{21}) + b_{22}\ Z_2(t-\tau_{22}),\end{aligned}\right\} \qquad (A.4)$$

Together with the initial conditions

$$Z_1(s) > |Y_1(s)|,\ \ s \in [-\tau, 0],\ i = 1,2. \qquad (A.5)$$

Then

$$|Y_1(t)| < Z_1(t) \text{ for } t \ge 0 \quad \text{and} \quad i = 1,2. \qquad (A.6)$$

PROOF. Let $U_1(t) = Z_i(t) - |Y_1(s)|$, $i = 1,2.$ $t \geq -\tau.$ Then we have from (A.2) and (A.4),

$$\left.\begin{aligned} \frac{dU_1(t)}{dt} &\geq -m_{11}\ U_1(t) + b_{11}U_1(t-\tau_{11}) + m_{12}\ U_2(t-\tau_{12}), \\ \frac{dU_2(t)}{dt} &\geq -m_{22}U_2(t) + m_{21}U_1(t-\tau_{21}) + b_{22}\ U_2(t-\tau_{22}), \end{aligned}\right\} t > 0, \quad (A.7)$$

With

$$U_1(s) = Z_i(s) - |Y_i(s)| > 0, \quad i = 1,2, s \in [-\tau, 0]. \quad (A.8)$$

Consider (A.7)-(A.8) first on the interval $[-\tau, 0]$ : It will follow from (A.7) that $U_1(t) \geq U_1(0) \exp[-m_{11}t]$

$$+\int_0^t \left[b_{11}U_1(s-\tau_{11}) + m_{12}\ U_2(s-\tau_{12})\right] \exp[-m_{11}(t-s)ds > 0,] \quad (A.9)$$

For all $t \in [0, \tau^*]$ since $U_i(s) > 0$ on $[-\tau^* 0]$ where $\tau^*$ = positive minimum of $(\tau_{11,}\tau_{12,}\tau_{21,}\tau_{22,})$. Similarly $U_2(t) > 0$ for $t \in [0, \tau^*]$. Now by a similar argument it will follow that $U_1(t) > 0, U_2(t) > 0$ for $t \in [\tau^*, 2\tau^*]$. Repeating the above procedure we obtain the result (A.6).

It follows from Lemmas 1 and 2 that in order to prove theorem 3.1 it is enough to prove that $Z_i(t) \to 0$ as $t \to \infty$ for $i = 1,2$ . which we will do as follows.

PROOF OF THEOREM 3.1 Define a Lyapunov-like function

$$H(Z_1,Z_2)(t)=Z_1(t)+Z_2(t)$$

$$+\int_{1-T}^t b_{11}\, Z_1(s)ds + \int_{1-T}^t m_{12}\, Z_2(s)ds$$

$$+\int_{1-T}^t m_{21}\, Z_1(s)ds + \int_{1-T}^t b_{22}\, Z_{22}(s)ds$$

Since each term of H in (A.10)is nonnegative for t>0, if we show that $H \to 0$ as $t \to \infty$ it will follow from Lemma 2 that

$$0 \leq |\, Y_{1,}(t)\,| \leq Z_1(t) \to 0 \text{ as } t \to \infty,$$

Implying that the perturbations $Y_{1,}$, $Y_2$, governed by (A.1)approach zero as $t \to \infty$, and this will prove our result Calculate the rate of change of H along the solutions of (A.4)so that

$$\frac{dH}{dt} = \frac{\partial H}{\partial t} + \frac{\partial H}{\partial Z_1}\frac{dZ_1}{dt} + \frac{\partial H}{\partial Z_2}\frac{dZ_2}{dt}$$

$$= -[Z_1(t)\{m_{11} - (b_{11} + m_{12})\} + Z_2(t)\{m_{22} - (b_{22}m_{12})\}] < 0$$

On using (2.10)since $Z_1(t), Z_2(t)$ are positive for all t>0. Thus $H(Z_1, Z_2)(t) \to 0$ as $t \to \infty$ since H is bounded below by zero: it follows from Lemma 6 below that $Z_1(t) \to 0, Z_2(t) \to 0$ as $t \to \infty$ the proof is complete.

The following set of results are intended to be a guide to prove result similar to the above in the case of continuously distributed delays An analogue of Lemma I is the following.

LEMMA 3 Let $K_{ij}(s)$,s∈[0,∞), $i,j = 1,2$,denote nonnegative continuous functions integrable on [0,∞) normalised such that

$\int_0^\infty sk_{ij}(s)ds< \infty$; $\int_0^\infty k_{ij}(s)ds=1$, I,j=1,2.

Consider the intergrodifferential system

$$\frac{dY_1(t)}{dt} = -m_{11}Y_1(t) + b_{11}\int_0^\infty k_{ij}(s)\, Y_1(t-s)ds$$

$$-m_{12}\int_0^\infty k_{12}(s)\, Y_2(t-s)ds$$

$$\frac{dY_2(t)}{dt} = -m_{22}Y_2(t) + m_{21}\int_0^\infty k_{21}(s)\, Y_1(t-s)ds$$

$$+b_{22}\int_0^\infty k_{22}(s)\, Y_2(t-s)ds$$

*Where* $m_{ij}, b_{ij}$ $(i,j = 1,2)$ *are positive constants Then*

$$\frac{d|Y_1(t)|}{dt} \leq -m_{11}|\,Y_1(t)\,| + b_{11}\int_0^\infty k_{11}(s)|Y_1(t-s)|ds,$$

$$+m_{12}\int_0^\infty k_{12}(s)|\,Y_2(t-s)|ds,$$

$$\frac{d|Y_2(t)|}{dt} \leq -m_{22}|\,Y_2(t)\,| + m_{21}\int_0^\infty k_{21}(s)|Y_1(t-s)|ds,$$

$$+b_{22}\int_0^\infty k_{22}(s)|\,Y_2(t-s)|ds,$$

*Where* $d|Y_1(t)||/dt$ *is defined by* (A.3).

Proof is again by direct verification.

The proof of the following lemma is similar to that of Lemma 2 above.

Lemma 4. *Consider (A.14) and the following intergrodifferential system:*

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1(t) + b_{11}\int_0^\infty k_{11}(s)Z_1(t-s)ds$$

$$+m_{12}\int_0^\infty k_{12}(s)Z_2(t-s)ds,$$

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1(t) + b_{11}\int_0^\infty k_{11}(s)Z_1(t-s)ds$$

$$+b_{22}\int_0^\infty k_{22}(s)Z_2(t-s)ds,$$

Together with the initial conditions

$$\infty > Z_i(s) > |Y_i(s)| for\ s \in (-\infty, 0).$$

(A.16)

Then $|Y_i(t)| < Z_i(t) for\ t \geq 0, i = 1,2.$

The following lemma will be useful to prove a result analogous to that of Theorem 3.1 corresponding to a class of continuously distributed time delays.

LEMMA 5. Let $Z_1(t), Z_2(t)$ be the solutions of (A.15)- (A.16) where $(Y_1\ Y_2)$ corresponds to an arbitrary but fixed solutions of (A.13). Then as a consequence of Lemma & it will follow that

$$0 \leq |Y_i(t)| \ Z_i(t) \to \infty \leq as\ t \to \infty, \qquad i = 1,2,$$

(A.17)

Whenever

$$\begin{cases} M_{11} > b_{11} + m_{21}, \\ M_{22} > b_{22} + m_{12} \end{cases}$$

PROOF. Define a Lyapunov –like function H where

$$H(Z_1,Z_2)(t) = Z_1(t) + Z_2(t) + b_{11}\int_0^\infty k_{11}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$$

$$+ m_{12}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$$

$$+ m_{21}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$$

$$+ b_{22}\int_0^\infty k_{22}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$$

Where $(Z_1,Z_2)$is any solution of (A.15)-(A.16) for arbitrary $(Y_1,Y_2)$defined on $(-\infty, 0]$ Computing the rate of change of H in (A.19)along the solution of (A.15)we derive that (on using (A.12))

$$\left(\frac{dH(Z_1Z_2(t)}{dt} = [Z_1(t)\{m_{11} - (b_{11} + m_{21})\} + Z_2(t)\{m_{22} - (b_{22} + m_{12})\}\right]$$

$< 0$(using (A.18)) for $Z_1(t) > 0\ and\ Z_2(t) > 0.$

By Lemma 6 below the result will follow

Lemma 6. Let H(t)=H$(Z_1, Z_2)(t)$ be as defined in (A.10) Suppose

$$m_{11} > b_{11} + m_{21} \text{ and } m_{22} > b_{22} |m_{12}$$

Then

$$Z_1(t) + Z_2(t) \to 0 \text{ast} \to \infty.$$

PROOF. Suppose (A.21) dose not hold; then there exists at least one sequence

$\{t_n\} \to \infty$ as n $\to \infty$ such that

$$[m_{11} - (b_{11} + m_{21})]\, Z_1(t_n) + [m_{22} - (b_{22} + m_{12})]Z_2(t_n) > \varepsilon \qquad (A.22)$$

For some sufficiently small positive number $\varepsilon$ and hence

$$\frac{dH(t_n)}{dt} < -\varepsilon, n = 1,2,3, \ldots.$$

By the continuity and differentiability of H with respect to t,

$$\mathrm{H}(t_n - \varepsilon) - \mathrm{H}(t_n) = \frac{dH(\mathrm{t}^*)}{dt}(-\varepsilon), \quad t_n - \varepsilon < t^* < t_n$$

If $\varepsilon$ is small we will have

$$= \frac{dH(\mathrm{t}^*)}{dt}\varepsilon/2$$

So that

$$\mathrm{H}(t_n) \leq \mathrm{H}(t_n - \varepsilon) - \varepsilon^2/2$$

$$\leq H(t_{n-1}) - \varepsilon^2/2$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2)$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2):$$

$$\leq H(t_0) - \mathrm{n}(\varepsilon^2/2): \mathrm{t}_0 = 0.$$

(A.12) and Lemma 5 will imply that $H(t_0) < \infty$ for any $t_0 > 0$. If n is sufficiently large it will follow that $H(t_n) < 0$ since $H(t_0) < \infty$, which contradicts the nonnegativity of H for all $t \geq 0$. Hence $Z_1(t) + Z_2(t) \to 0$ as $t \to \infty$ implying the asymptotic stability of the trivial solution of (A.1)(since $Z_1$ dominates $y_1$)

We note the result of Lemma 5 and 6 will remain valid if (A.18) is replaced by the condition $\mu(A) + \|B\| < 0$

Where A and B are defined as in (6.3).

We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities,

apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems

**Gopalsamy**

# Model For Innocuously Benign Delay In A Periodic System

**Abstract**

Sufficient conditions are obtained for the existence of a unique linearly denomination stable positive periodic solution of a system model of two competitions in a periodic environment with time lags in interspecific interactions. It is shown that if the self- regulating intraspecific interaction effects are strong enough and act without time delays then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a periodic domination. An example in finance will make things clear. Corporate finance for one big plant generates lot of accounts of the workers whose salary shall be disbursed to the debit of cash credit account. Here Cash credit account is a resource as also the one to which debits are raised, a la predator. Some of the accounts from which the credits come are demand bills purchase, usance bills discount, vostro and nostro accounts in respect of foreign bills and Letters of credit. In economics inventories are used to produce output, output is sold to produce income which might be saved in part and recycled in investment. Cosmologically take an example of a Black hole which gormandizes matter and radiates energy. Notwithstanding information paradox, one thing is clear that matter is being created in to radiation. We show how the model could be applied to axiomatic predications, postulation alcovishness, phenomenological methodologies, propositional statement, and other problems in cosmology, philosophy, and even in Green House effects. We study the sine qua non properties of such a system. In fact these models can be applied to such unsolved problems like Yang mills Theory and mass gap, and other problems. As we apply the model, we shall go on dilating the position status and the applicability of the model. We shall not touch the electrical engineering problems which might rise questions like what is charge of a Bank. We may say when a customer enters a Bank; we can say the Bank is charged. These factors are still relativistic and we shall make an in-depth study as we progress. Gopalsamy's research was published in Australian Journal of Mathematical Society vide: B 25349=365 It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 1. Introduction

At the outset we assume that the system satisfies Haimovici is caution. The purpose of this article is to establish a set of sufficient conditions for the existence of a unique linearly asymptotically stable positive periodic solution for a nonautonomous system of the form

$$\frac{dx(t)}{dt} = x(t)f_1\big(t, x(t), y(t-\tau_1)\big),$$

**(1.1)**

$$\frac{dy(t)}{dt} = y(t)f_2\,\big(t, x(t-\tau_2), y(t)\big).$$

In the context of confederation dynamics a model system of the form (1.1) means that the interspecific interactions occur with time delays $\tau_1$ 0, $\tau_2$ 0 in a periodic environment, if $f_1$ *and* $f_2$ are assumed to be periodic in $t$ with a common period $\omega$ 0. Note the interactions are inter specific.

Although model systems in periodic domains have been investigated by a large number of authors (for some references see Gopalsamy (4], [5]), periodically varying environments in time delayed model systems of two or more interacting species have not been investigated. The need for Multitype system models in periodic environments and time delayed interactions are evident from the extensive literature of that type relating to single species. While it is known that periodic ecosystems can support more species than uniform environments (Gopalsamy [4]), time delays might destabilize the system leading to a possible loss of one or more species. Furthermore, periodic environments in population dynamics have some relevance to natural selection and evolution; we will consider this aspect briefly at the end.

The principal finding of the following analysis of (1.1) is that, under suitable hypotheses on $f_1$ *and* $f_2$, the system (1.1) has a linearly asymptotically stable unique positive periodic solution for all nonnegative delays $\tau_1, \tau_2$ (while our numerical experiments suggest that the periodic solution of (1.1) is globally asymptotically stable and is independent of the delay parameters $\tau_1, \tau_2$ it is not known how one can prove such a strong result). Since time delays are usually associated with possible destabilization of otherwise stable systems, we can describe the time delays appearing on (1.1) as “harmless”.

It has been recently shown by the author (Gopalsamy [6]) that if the self –regulating negative feedback in each denomination is strong enough and acts without time delays, then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a temporally uniform domain. The result of this paper is thus an extension to a periodic environment of the "harmless" nature of interspecific interaction delays in a temporally constant environment mentioned above.

## 2. Existence and uniqueness of periodic solutions

We will now formulate the problem of existence of periodic solutions of the periodic delay differential system

$$\frac{dx_1(t)}{dt} = x_1(t) f_1\big(t, x_1(t), x_2(t-\tau_{12})\big),$$

$$t > 0, \qquad \textbf{(2.1)}$$

$$\frac{dx_2(t)}{dt} = x_2(t) f_2\big(t, x_1(t-\tau_{21}), x_2(t)\big),$$

(Where $\tau_{12}, \tau_{21}$ are nonnegative constants) in a form suitable for the application of fixed point theory of positive operators. We make the following assumptions regarding $f_1$ and $f_2$ in (2.1).

($H_1$) The partial (Frechet) derivatives of $f_1, f_2$ satisfy the following:

$$\frac{\partial f_i}{\partial x_j} < 0, \qquad i, j = 1, 2 \text{ For all } t > -\tau \text{ where } \tau = \max(\tau_{12}, \tau_{21}). \qquad \textbf{(2.2)}$$

($H_2$) $f_1(t, x_1, x_2) \equiv f_1(t + \omega, x_1, x_2)$ for all $t > -\tau, x_1 \geq 0, x_2 \geq 0$. **(2.3)**

($H_3$) There exist positive constants $b_{il}$, $b_i^u$, $a_{il}$, $a_i^u$ ($l = 1, 2$) such that

$$b_{il} - a_{il} x_i < f_i(t, x_1, x_2) < b_i^u - a_i^u x_i,$$

$$b_{il}/a_{il} < b_i^u/a_i^u, \qquad \textbf{(2.4)}$$

For $i$= 1, 2 and $(x_1, x_2) \in R_+ \times R_+$. If $\tau = 0$ then one can show that (2.1)-(2.4) will imply that the nondelayed periodic system (2.1) will have an asymptotically stable strictly positive periodic solution (for details see Gopalsamy [5]); this result will become a special case of what we obtain in the following.

We will use the following notation:

$$B_\omega^\tau = \{h: [-\tau, \infty) \to R, h \text{ is continuous}, h(t + \omega) = h(t), t \geq -\tau\},$$

$$B = B_\omega^\tau \times B_\omega^\tau,$$

$$K_1 = \{ h \in B_\omega^\tau | h(t) \geq 0, t \in [-\tau, \infty)\},$$

$$K = K_1 \times K_1,$$

$$||h|| = \left\{ \sup_{t \in [0,\omega],} | h(t)|, h \in B_\omega^\tau \right\}$$

$$||x_1, x_2|| = \max\{||x_1||, ||x_2||, (x_1, x_2) \in B\},$$

$$G_1 = \{(x_1, x_2) \in K, ||(x_1, x_2)|| \leq b_l = \min(b_{1l} / a_{1l}, b_{2l}/a_{2l})\}$$

$$G_2 = \left\{(x_1, x_2) \in K, ||(x_1, x_2)|| \leq b^u = \max\left(\frac{b_1^u}{a_1^u}, \frac{b_2^u}{a_2^u}\right)\right\},$$

$\bar{G} =$ closure of $G_2$ *in* $B$ and

$\frac{\bar{G}_2}{G_1} =$ complement of $G_1 in\ \bar{G}_2;\ G_1 \subset G_2$ by **(2.4)**

The cone $K = K_1 \times K_1$ induces a partial order in $B$ and we will write $(x_1, x_2) \leq (y_1, y_2)$ if and only if $x_1 - y_1 \in K_1$. we define an operator $N$ on $B$ as follows

$$N(x_1, x_2)(t) = \begin{matrix} N_1(t, x_1, x_2) \\ N_2(t, x_1, x_2) \end{matrix} \quad \text{for t> } 0, \qquad \textbf{(2.5)}$$

Where

$$N_1(t, x_1, x_2) = \int_0^\omega G_{11}(m, \omega, t-s) g_1\left(m, s, x_1(x) x_2(s - \tau_{12})\right) ds$$

$$+ \int_0^1 G_{12}(m, 0, t-s) g_1\left(m, s, x_1(s), x_2(s - \tau_{12})\right) ds,$$

$$N_2(t, x_1, x_2) = \int_0^\omega G_{22}(m, \omega, t-s) g_2\left(m, s, x_1(s - \tau_{21}), x_2(s)\right) ds \qquad \textbf{(2.6)}$$

$$+ \int_0^1 G_{21}(m, 0, t-s) g_2\left(m, s, x_1(s - \tau_{21}), x_2(s)\right) ds,$$

For $t > 0$, with

$$g_1\left(m, s, x_1(s), x_2(s - \tau_{12})\right) = m x_1(s) + x_1(s) f_1\left(s, x_1(s), x_2(s - \tau_{12})\right),$$

**(2.7)**

$$g_2\big(m, s, x_1(s-\tau_{21}), x_2(s)\big) = mx_2(s) + x_2(s) f_2\big(s, x_1(s-\tau_{21}) x_2, (s)\big),$$

$$G_{ll}(m, \omega, t-s) = e^{-m\omega}(1-e^{-m\omega})^{-1} e^{-m(t-s)}, \qquad i = 1, 2,$$

**(2.8)**

$$G_{ij}(m, 0, t-s) = e^{-m(t-s)}, i \neq j, i, j = 1, 2,$$

$$N(x_1, x_2)(t) = \begin{cases} N_1(n\omega + t, x_1, x_2) \\ N_2(n\omega + t, x_1, x_2) \end{cases} for\ t \in [-\tau, 0], \qquad \textbf{(2.9)}$$

*m* being a suitably selected positive constant (more will be said below) while *n* is a positive constant such that $n\omega > \tau$.

We note immediately the following from the definition of $N$: $if\ (x_1, x_2) \in \bar{G}/G_1$ then $N(x_1, x_2)(t)$ ☐ and if $N(x_1, x_2)$ is periodic $\omega$ for $t$ ☐ then $N(x_1, x_2)$ is periodic in *t* with period $\omega$ for $t$ ☐$-\tau$. The existence of at least one positive periodic solution of (2.1)-(2.4) is established by means of a few lemmas and a result of Schmitt [9].

Lemma 1. *A necessary and sufficient condition for* $(x_1(t), x_2(t))$ , $t$ ☐$-\tau$ *to be a periodic solution of (2.1)-(2.4) of period* $\omega$ *is that* $(x_1, x_2)$ *be a fixed point of N such that*

$$N(x_1, x_2)(t) = \big(x_1(t), x_2(t)\big) \quad and\ (x_1, x_2) \in K \cap \left(\frac{\bar{G}}{G_1}\right).$$

PROOF. The details of proof are similar to those of Lemma 2 of Busenberg and Cooke [1]. We first verify the continuity of $N(x_1, x_2)$ assuming that $N(x_1, x_2)(t)$ is periodic in *t* of period $\omega$. The continuity of $N(x_1, x_2)$ for $t > 0$ follows from the definition in (2.5)-(2.8) WHICH ALSO PROVIDES $\lim_{t\to 0+} N(x_1, x_2)(t)$ . It will follow from (2.9) and the assumed periodicity of $N(x_1, x_2)\ that\ N(x_1, x_2)$ is continuous

For $t$ ☐$-\tau$. Now we will verify the periodicity of $N(x_1, x_2)$ in $t$. A simple computation leads to

$$N_1(t+w, x_1, x_2) - N_1(t, x_1, x_2) = e^{-mt} F_1(t),$$

**(2.10)**

Where

$$F_1(t) = \int_0^{\omega} G_{11}\,(m, \omega, t-s\,)g_1\big(m, s, x_1(s), x_2(s-\tau_{12})\big)ds$$

$$+\int_0^{\omega} G_{12}(m, 0, t-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$+\int_{\omega}^{\omega+t} G_{12}(m, 0, \omega-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$-\int_0^{\omega} G_{11}(m, \omega-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$+\int_0^{t} G_{12}(m, 0, -s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

**(2.11)**

Which on simplification using the periodicity of $f_1$ shows that $F_1(t) \equiv 0$ whenever $x_1, x_2, \in B$, thus implying the periodicity of $N_1$ in $t$ of period$\omega$. A similar argument for $N_2$ proves that $(x_1, x_2) \in B \Rightarrow N\,(x_1, x_2) \in B$.

Now suppose $(x_1, x_2)$ is a fixed point of *N and let* $(x_1, x_2)\ \in K \cap (\bar{G}/G_1\,)$. Then it is easily seen that $(x_1, x_2)$ is a solution of (2.1) and since $(x_1, x_2)\ \in K \cap (\bar{G}/G_1\,)$ it will follow that $(x_1, x_2)$ is a periodic solution of (2.1).

Conversely, let $(x_1, x_2)$ be a periodic solution of period $\omega$ of (2.1). Then $(x_1, x_2)$ is necessarily a periodic solution of

$$\frac{dx_1(t)}{dt} + mx_1(t) = g_1\,\big(m, t, x_1\,(t), x_2(t-\tau_{12})\big)$$

**(2.12)**

$$\frac{dx_2(t)}{dt} + mx_2(t) = g_2\,\big(m, t, x_1\,(t-\tau_{21}), x_2(t)\big),$$

For any arbitrary real constant *m* and hence for any fixed constant $m > 0$. It follows from (2.12) that any periodic solution of (2.1) and hence of (2.12) is a fixed point of $N(x_1, x_2)$. We claim that such a fixed point is in $K \cap (\bar{G}_2/G_1)$. Suppose $(x_1, x_2) \notin K \cap (\bar{G}_2/G_1)$. the following possibilities can arise:

(i) Either $x_1, x_2$ or both $x_1$ *and* $x_2$ may become negative for some value of *t* and it will follow from (2.1) that subsequent value of $x_1, x_2$ or $x_1$ *and* $x_2$ will be negative and hence from (2.4) and (2.1) their derivatives will be negative for all subsequent *t* , implying that $x_1, x_2$ or $x_1$ *and* $x_2$ cannot be periodic.

(ii) While $(x_1, x_2) \in K \subset B$ there may be a value of *t*, say $t^*$, such that $x_1(t^*) > b_1^u/a_1^u$ or $x_2(t^*) > b_1^u/a_1^u$ and in such a case $x_1(t) or\ x_2(t)$ cannot be periodic since respectively $dx_1(t^*)/dt < 0\ or\ dx_2(t^*)/dt < 0,$ which will imply that the values of $x_1(t^*)\ or\ x_2(t^*)$ cannot be repeated from below.

(iii) Similarly if there is a $t^*$ such that $0 < x_1(t^*) < b_{1l}/a_{1l}\ or\ 0 < x_2(t^*) < b_{2l}/a_{2l}$ then such solutions cannot be periodic since $dx_1(t^*)dt > 0$ or $dx_2(t^*)dt > 0$, which will mean that the values of $x_1(t^*)$ or $x_2(t^*)$ cannot be repeted from above.

Thus every periodic solution of (2.1)-(2.4) is a fixed point of *N* in $K \cap \left(\frac{G_2}{G_1}\right) \subset B$.

**LEMMA 2.** *Let N be as in (2.5)-(2.6). Suppose* $\lambda \in (0,1) and\ \lambda N(x_1, x_2) = (x_1, x_2)$ *for some* $((x_1, x_2) \in K - \{0\}$. *then there exist* $\alpha_1^* > 0, \alpha_2^* > 0$ *such that*

$$0 < x_1(t) \leq \alpha_1^* < b_1^u / a_1^u,$$

$$0 < x_2(t) \leq \alpha_2^* < b_2^u / a_2^u,$$

*for* $t \in [-\tau, \infty)$.

PROOF. If $\lambda N(x_1, x_2) = (x_1, x_2)$ then from the definition of N we get for all $t > 0$ and $(x_1, x_2) \in K - \{0\}$,

$$\frac{dx_1(t)}{dt} = \lambda\, x_1\,(t) f_2(t, x_1(t), x_2\,(t-\tau_{12})) + x_1\,(t) m(\lambda - 1)$$

**(2.13)**

$$\frac{dx_2(t)}{dt} = \lambda\, x_2\,(t) f_2(t, x_1(t-\tau_{21}), x_2\,(t)) + x_2\,(t) m(\lambda - 1)$$

For all $t$ such that $x_1(t) \geq b_1^u/a_1^u$ *or* $x_2(t) \geq b_2^u/a_2^u$ we have $\frac{dx_1}{dt} < 0$ *or* $\frac{dx_2}{dt} < 0$; but these are impossible since $x_1, x_2$ are periodic in t. Thus by the continuity and periodicity of $x_1, x_2$ we have $x_1(t) \leq x_1^* < b_1^u/a_1^u$, $x_2(t) \leq x_2^* < b_2^u/a_2^u$ for all $t \geq 0$; by the definition and periodicity the same will follow for all $t \geq -\tau$. A similar argument holds when $x_1(t) < 0$ or $x_2(t) < 0$, noting that in such a case $dx_1/dt > 0$ or $dx_2/dt > 0$, again contradicting the periodicity of $(x_1, x_2)$ in *t*.

LEMMA 3. *The operator N maps* $K \cap (G_2/G_1)$ *in to K and is completely continuous*

PROOF. From the definition of N and (2.4) we derive that

$$N_1(t, x_1, x_2) \geq \int_0^\omega G_{11}\,(m, \omega, t-s) x_1(s)[m + b_{1l} - a_{1l} x_1(s)]ds$$

$$+ \int_0^\omega G_{12}\,(m, 0, t-s) x_1(s)[m + b_{1l} - a_{1l} x_1(s)]ds.$$

**(2.14)**

For $(x_1, x_2) \in K \cap (G_2/G_1)$, we have $0 < x_1 < b_1^u/a_1^u$ and hence we get $N_1(t, x_1, x_2) > 0$ for all $t \geq 0$ by a proper choice of the positive constant say $m_1$.

We make a similar choice say $m_2$ for *m* for the case $N_1(t, x_1, x_2) \geq 0$ and then choose *m* = max $(m_1, m_2)$. Thus we have

$$N: K \cap (G_2/G_1) \to K \qquad for\ t \geq 0,$$

And by the periodicity of N we have the same for $t \to -\tau$.

To prove the complete continuity of N we proceed as follows. We know that if $(x_1, x_2) \in B$ then $N(x_1, x_2) \in B$ and so we can restrict $t$ to the set $[0, n\,\omega], n\omega > r$. We write $N = L \circ Q$ where $L$ is the linear operator, $L : B \to B$, $L(u_1, u_2)(t)$

$$= \int_0^{\omega} G_{11}(m, \omega, t-s) \begin{bmatrix} u_1 & (s) \\ u_2 & (s) \end{bmatrix} ds, + \int_0^{t} G_{12}(m, 0, t-s) \begin{bmatrix} u_1 & (s) \\ u_2 & (s) \end{bmatrix} ds$$

**(2.15)**

And Q is the nonlinear operator

$$Q(u_1, u_2)(t) = \begin{matrix} g_1(m, t, u_1(t), u(t-\tau_{12})) \\ g_2(m, t, u_1(t-\tau_{21}), u_2(t)) \end{matrix} for\ \ t > 0,$$

**(2.16)**

$$Q(u_1, u_2)(t) = Q\,(u_1, u_2)(n\omega + t) \quad for\ t \in [-\tau, 0].$$

It can be found that

$$||L(u_1, u_2)(t)|| \leq max \left\{ \frac{m+b_1^u}{m}, \frac{m+b_2^u}{m} \right\} ||u_1, u_2||,$$

**(2.17)**

$$\left\| \frac{d}{dt} L(u_1, u_2) \right\| \leq max \left\{ \frac{m+b_1^u}{m}, \frac{m+b_2^u}{m} \right\} ||u_1, u_2||.$$

The complete continuity of the linear map $L$ follows from (2.17). Now since Q is continuous and maps bounded sets into bounded sets it will follow that $L \circ Q$ is completely continuous.

To prove the existence of periodic solution of (2.1)-(2.4) we will use the following result of Schmitt [9].

**THEOREM (Schmitt).** *Let $H_1$ and $H_2$ be nonempty bounded open neighborhoods of the origin in a Banach space B with a positive cone K such that $H_1 \subset H_2$ and let N: $K \cap (H_2 / H_1) \to K$ be completely continuous and which satisfies the following conditions:*

***(a)*** *There exists $(z_1, z_2) \in K \subset B$, $\|(z_1, z_2)\| = 1$ and a $\mu$ satisfying*

$$\mu > sup\{\|(x_1, x_2) - N(y_1, y_2)\|, (x_1, x_2) \in K \cap H_1, (y_1, y_2) \in K \cap \partial H_2\}$$ *such that all solution* $(x_1, x_2) \in K \cap (H_2/H_1)$ *of* $(x_1, x_2) = N(x_1, x_2) + \lambda(z_1, z_2) with\ 0 < \lambda < \mu\ satisfy\ (x_1, x_2) \notin \partial H_1.$

*(b) All solution* $(x_1, x_2) \in K \cap (H_2/H_1)\ of\ (x_1, x_2) = \lambda N(x_1, x_2), 0 < \lambda < 1\ satisfy\ (x_1, x_2) \notin \partial H_2.$

*Then N has at least one fixed point in* $K \cap (H_2/H_1)$.

We can now present our existence result.

**THEOREM 1.** *Assume that* $f_1, f_2$, *of* (2.1) *satisfy* (2.2)-(2.4). *then there exists at least one positive periodic solution of* (2.1)-(2.4).

PROOF. By Lemma 1, it is enough to show that the operator $N: K \cap (G_2/G_1) \to K \subset B$ has at least one fixed point in $G_2/G_1$ and we will do this by applying the result of Schmitt [9]. First we let

$$H_2 = \{(x_1, x_2) \in B | \ ||x_1, x_2|| < \max(b_1^u/a_1^u, b_2^u/a_2^u)\} \equiv G_2,$$

$$H_1 = \{(x_1, x_2) \in B | \ ||x_1, x_2|| < \min(b_{1l}/a_{1l}, b_{1l}/a_{1l})\} \equiv G_1.$$

By Lemma 3, $N: K \cap (H_2/H_1) \to K$ is completely continuous. Choose $\mu = \infty$ and $(z_1, z_2) \equiv (1.1)$. Then all solution $(x_1, x_2) \in K \cap (H_2/H_1)$ of $(x_1, x_2) = N(x_1, x_2) + \lambda(z_1, z_2), 0 < \lambda < \infty$, satisfy (2.1) and hence by their periodicity $(x_1, x_2) \notin \partial H_1$. The condition (a) of Schmitt's theorem is satisfied since $\mu = \infty$ by choice and $(x_1, x_2)$ and $(z_1, z_2)$ are periodic. By Lemma 2 the condition (b) of Schmitt's theorem is satisfied. Thus there exists at least

one fixed point of N in $K \cap (H_2/H_1)$ and hence there exists at least one positive periodic solution of (2.1)-(2.4) and the proof is complete.

We remark that so far no restriction on the magnitudes of the delays $\tau_{12}, \tau_{21}$ has been imposed except that they be nonnegative.

## 3. Uniqueness of the periodic solution

Applications of fixed point theorems to prove the existence of solutions of nonlinear problems usually suffer from the drawback that the uniqueness of solutions has to be proved often by other methods. We will use the following result of Smith [11] which is a slight variation of his Theorem 3.5 (see also Busenberg and Cooke [1]).

THEOREM (Smith). *Let N be an operator on a cone K in the Banach space B and let* $K^* = \{(x_1, x_2) \in K \cap B, 0 \leq$ *x1,x2≤C∗ where C∗some positive constant is. Assume the following:*

(a) If $(x_1, x_2) \in K^* - \{(0,0)\}$ and $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

(b) N is monotonic nondecreasing on $K^*$.

(c) If $(x_1, x_2)$ is in the interior of $K^*$ and if $0 < \lambda < 1$ then there exists an $\eta > 0$ depending on $(x_1, x_2)$ and $\lambda$ such that $N(\lambda x_1, \lambda x_2) \geq \lambda(1+\eta) N(x_1, x_2)$. Then N has at most one fixed point different from $(0,0)$ in $K^* \subset B$.

The cone K= $K_1 \cap K_1$ where $K_1 = \{h \in B_\omega^\tau : h(t) \geq 0, t \in [-\tau, \infty)\}$ is slightly too wide in that solutions of the form $\{x_1(t), 0\}$ and $\{0, x_2(t)\}$ for (2.1) are also possible. We choose a narrower cone $K^*$ in $B$ as follows: define

$$S = \{(x_1, x_2) \in B : b_{1l}/a_{1l} \leq \|x_1\| \leq b_1^u/a_1^u,\ b_{2l}/a_{2l} \leq \|x_2\| \leq b_2^u/a_2^u,\}, \quad \textbf{(3.1)}$$

And let

$$K^* = \{\lambda(x_1, x_2) \in B : (x_1, x_2) \in S, \lambda \geq 0\}. \quad \textbf{(3.2)}$$

Since $\hat{G}_2/G_1 \subset K^*$, it is not difficult to see that N: $\hat{G}_2/G_1 \to B$ has at least one fixed point which is a periodic solution of (2.1)-(2.4). We will prove that this solution is unique by applying the result of Smith cited above. The following lemmas are needed first.

LEMMA 4. If $(x_1, x_2) \in K^* - \{0.0\}$ and if $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

PROFF. If $(x_1, x_2) = N(x_1, x_2)$ and $(x_1, x_2) \in K^* - \{0,0\}$, then $(x_1, x_2)$ is a periodic solution of (2.1) –(2.4) and hence we have

$$(b_{1l} - a_{1l}x_1)x_1 < \frac{dx_1}{dt} < (b_1^u - a_1^u x_1)x_1,$$
$$(b_{2l} - a_{2l}x_2)x_2 < \frac{dx_2}{dt} < (b_2^u - a_2^u x_2)x_2, \qquad (\mathbf{3.3})$$

And hence any nonzero periodic solution $(x_1, x_2)$ of (2.1)-(2.4) has to be such that

$$b_{1l}/a_{1l} < x_1(t) < b_1^u - a_1^u,$$

$$b_{2l}/a_{2l} < x_2(t) < b_2^u - a_2^u,$$

For $t > -\tau$, which proves the lemma.

LEMMA 5. The mapping $N$ defined by (2.5)-(2.9) is monotone nondecreasing on $K^*$.

PROOF. Let $(x_1, x_2) \in K^*$, $(y_1, y_2) \in K^*$ such that $(x_1 - y_1, x_2 - y_2) \in K^*$ and $(x_1, x_2) \neq (y_1, y_2)$. Then $x_1 \geq y_1$, $x_2 \geq y_2$. From the definition of $N$ we derive that

$N_1(t, x_1, x_2) - N_1(t, y_1, y_2)$

$$= \int_0^\omega G_{11}(m, \omega, t-s)\{m[x_1(s)] + x_1(s)f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s)f_1(s, y_1(s), y_2(s-\tau_{12}))\}ds$$

$$+ \int_0^t G_{12}(m, 0, t-s)$$

$$x\{m[x_1(S) - y_1(s)] + [x_1(s)f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s)f_1(s, y_1(s), y_2(s-\tau_{12}))]\}ds \qquad \mathbf{(3.4)}$$

With a similar relation for $N_2(t, x_1, x_2) - N_2(t, y_1, y_2)$. Since $f_1$ is continuous, $f$ is bounded on $K^*$; also it will follow from $(x_1, x_2) \neq (y_1, y_2)$ that

$$0 \not\equiv x_1(s) - y_1(s) \geq 0 \text{ and } 0 \not\equiv x_2(s) - y_2(s) \geq 0;$$

A sufficiently large positive constant $m$ will make the right side of (3.4) positive.A similar argument will lead to $N_2(t, x_1, x_2) - N_2(t, y_1, y_2) > 0$ perhaps with a larger constant $m$. Finally we can choose the bigger of these two constants $m$ to yield

$$N(x_1, x_2) - N(y_1, y_2) \text{ for } (x_1 - y_1, x_2 - y_2) \in K^* - \{(0,0)\}.$$

The monotonicity of $N$ on $K^*$ follows.

LEMMA 6. Let $(x_1, x_2) \in$ interior of $K^*$ and let $\lambda$ be $a$ constant such that $0 < \lambda < 1$. Suppose $f_1, f_2$ satisfy (2.2) - (2.4). Then there exists $a$ Positive constant $\eta$ such that

$$N(\lambda x_1, \lambda x_2) \geq \lambda(1+\eta) N(x_1, x_2). \quad \textbf{(3.5)}$$

PROOF. A direct computation leads to

$$N_1(t, \lambda x_1, \lambda x_2) - \lambda N_1(t, x_1, x_2)$$

$$= \lambda \int_0^{\omega} G_{11}(m, \omega, t-s) g_1\big(s, x_1(s), x_2(s-\tau_{12})\big) H_1(s, \lambda, x_1, x_2) ds$$

$$+ \lambda \int_0^{t} G_{11}(m, 0, t-s) g_1\big(s, x_1(s), x_2(s-\tau_{12})\big) H_1(s, \lambda, x_1, x_2) ds$$

**(3.6)**

Where

$$H_1(s, \lambda, x_1, x_2) = \frac{m\, x_1(s) + x_1(s) f_1\big(s, \lambda x_1(s), \lambda x_2(s-\tau_{12})\big)}{m\, x_1(s) + x_1(s) f_1\big(s, x_1(s), x_2(s-\tau_{12})\big)} - 1. \quad (\mathbf{3.7})$$

Since $f_1$ is decreasing (see (2.2)) with respect to $x_1$ and $x_2$ there exists a positive constant say $\eta_1$ such that

$$H_1(s, \lambda, x_1, x_2) \geq \eta_1 \text{ for } 0 < \lambda < 1, s > 0, (x_1, x_2) \in K^*, \quad \textbf{(3.8)}$$

And hence

$N_1(t, \lambda x_1, \lambda x_2) \geq \lambda(1+\eta_1) N_1(t, x_1, x_2)$.

A similar argument will show that there exists an $\eta_2 > 0$ such that

$N_2(t, \lambda x_1, \lambda x_2) \geq \lambda(1+\eta_2) N_2(t, x_1, x_2)$,

And the result follows if we choose $\eta = \min(\eta_1, \eta_2)$.

Now the following uniqueness result is a direct consequence of the result of Smith [11] and the Lemmas 4, 5, 6.

THEOREM 2. *Suppose $f_1, f_2$ of (2.1) satisfy (2.2) – (2.4). Then there exists a unique $\omega$ - periodic strictly positive periodic solution of (2.1).*

## 4. Linear stability of the periodic solution

Let $\{x_1^*(t), x_2^*(t), t \geq -\tau\}$ be the strictly positive unique periodic solution of (2.1). Let $\{x_1(t), x_2(t), t \geq -\tau\}$ be any other solution of (2.1). To study the linear stability of $(x_1^*, x_2^*)$ we let

$$\left.\begin{aligned} x_1(t) &\equiv x_1^*(t)\{1 + X_1(t)\}, \\ x_2(t) &\equiv x_2^*(t)\{1 + X_2(t)\}, \end{aligned}\right\} \quad t \geq -\tau, \tag{4.1}$$

And derive that the perturbations $(x_1, x_2)$ are governed by the linear variational system (after neglecting the non linear terms in $x_1, x_2$)

$$\left.\begin{aligned} \frac{dX_1(t)}{dt} &= -a_{11}(t)\, x_1^*(t) X_1(t) - a_{12}(t) x_2^*(t-\tau_{12})\, X_2(t-\tau_{12}) \\ \frac{dX_2(t)}{dt} &= -a_{21}(t) x_1^*(t-\tau_{21}) x_1(t-\tau_{21}) - a_{22}(t) x_2^*(t) x_2(t), \end{aligned}\right\} \tag{4.2}$$

For $t > 0$, where

$$\left.\begin{aligned} -a_{11}(t) &= \frac{\partial f_1}{\partial x_1} \\ -a_{12}(t) &= \frac{\partial f_1}{\partial x_2} \end{aligned}\right\} \text{Evaluated at } (t, x_1^*(t), x_2^*(t-\tau_{12})), \tag{4.3}$$

$$\left.\begin{aligned} -a_{22}(t) &= \frac{\partial f_2}{\partial x_1} \\ -a_{22}(t) &= \frac{\partial f_2}{\partial x_2} \end{aligned}\right\}$$ Evaluated at $(t, x_1^*(t-\tau_{21}), x_2^*(t))$.

To establish the linear asymptotic stability of $(x_1^*(t), x_2^*(t))$ it is enough to prove that of the trivial solution of (4.2) - (4.3) We will do this by means of a combination of a comparison differential system and a Lyapunov function. First we define

$$\frac{d|X_1(t)|}{dt} = \begin{cases} \frac{dX_1(t)}{dt} & if\ X_i(t) \geq 0, \\ -\frac{d\setminus X_1(t)}{dt} & if\ X_i(t) < 0, \end{cases} \quad i = 1,2, \qquad \textbf{(4.4)}$$

And derive from (4.2)- (4.3) the following comparison system of differential inequalities:

$$\begin{aligned} \frac{d|X_1(t)|}{dt} &\leq -b_{11}, x_1^*(t)\,|X_1(t)| + b_{12}x_2^*(t-\tau_{12})\,|X_2(t-\tau_{12})|, \\ \frac{d|X_2(t)|}{dt} &\leq -b_{21}, x_1^*(t-\tau_{21})\,|X_1(t-\tau_{12})| - b_{22}x_2^*(t)\,|X_2(t)|, \end{aligned} \qquad \textbf{(4.5)}$$

For $> 0$, where

$$-b_{11} = \max_{t\in-\tau}\{-a_{11}(t)\},$$

$$\begin{aligned} b_{12} &= |\min_{t\in-\tau}\{-a_{12}(t)\}| \\ b_{21} &= |\min_{t\in-\tau}\{-a_{21}(t)\}| \end{aligned} \qquad \textbf{(4.6)}$$

$$-b_{22} = \max_{t\in-\tau}\{-a_{22}(t)\}$$

We can now formulate the result as follows. By the periodicity of $a_{ij}$, the constants $b_{ij}$ are independent of $\tau$ (see also (2.2)).

THEOREM 3. *The unique strictly positive periodic solution* $\{x_1^*(t), x_2^*(t), t\in-\tau\}$ of $(2.1)-(2.4)$ *is linearly asymptotically stable provided*

$$|-b_{11}| > b_{21} \text{ and } |-b_{22}| > b_{12},$$

**(4.7)**

*Where the constants* $b_{ij}$ $(i, j = 1,2)$ *are as defined in* (4.6).

PROOF. Consider the scalar valued Lyapunov-like function $v(t) = V(t, X_{1,}X_{2,})$where

$$v(t) = V(t, X_{1,}X_{2,}) = |X_{1,}(t)| + |X_{2,}(t)| \int_{t-\tau_{12}}^{t} b_{12} x_2^*(u)|X_{2,}(u)|du$$

$$+\int_{t-\tau_{21}}^{t} b_{21} x_1^*(u)|X_{1,}(u)|du. \quad \textbf{(4.8)}$$

Since $x_1^*, x_2^*$ are strictly positive, it is found from (4.7)-(4.) that

$$\frac{dv}{dt} \leq -[(|b_{11}| - b_{21})\, x_1^*(t)|X_{1,}(t)| + (|b_{22}| - b_{22})\, x_2^*(t)|X_2(t)|] \quad \textbf{(4.9)}$$

For $t > 0$. It follows from (4.7)-(4.9) that $V(t, X_{1,}X_{2,}) \geq |X_{1,}(t)| + X_{2,}(t)|$ and $dv/dt \leq 0$,which will imply the stability of the trivial solution of (4.2)-(4.3). We claim that (4.9) also implies the asymptotic stability of the trivial solution of (4.2)-(4.3)in the sense that

$$\lim_{t-\infty} [\,|X_{1,}(t)| + |X_{2,}(t)|\,] = 0,$$

**(4.10)**

For all bounded initial conditions $X_{1,}(s), X_{2,}(s), s \in \{\mathcal{T}, 0\}$. $Suppose (4.10) is\ not$ true Then we can find an increasing sequence $\{t_m\} \to \infty$ as $m \to \infty$ such that

$$(|b_{11}| - b_{21})\, x_1^*(t_m)|X_{1,}(t_m)| + (|b_{22}| - b_{12} x_2^*(t_m)|X_2(t_m)| > \varepsilon,$$

M=1, 2, 3…

**(4.11)**

Using the right derivative *dv/dt* of $v$ we derive that

$$v(t_m - \varepsilon) - \mathrm{v}(t_m) \leq -\int_{t_m-\varepsilon}^{t_m} \frac{\mathrm{dv(T)}}{\mathrm{dT}}\,\mathrm{d}\tau.$$

**(4.12)**

If $\varepsilon$ is small enough we have from (4.9)(4.11),(4.12)that $dv(\tau)/dr < (-\varepsilon/2)$ and

$$v(t_m - \varepsilon) - \mathrm{v}(t_m) \geq \varepsilon^2/2 \text{ and hence}$$

$$v(t_m) \geq v(t_m - \varepsilon) - (\varepsilon^2/2)$$

$$\geq v(t_m - 1)\text{-}(\varepsilon^2/2)$$

$$\geq v(t_m - 2)\text{-}(\varepsilon^2/2)$$

$$\geq v(t_0)\text{-m}\,(\varepsilon^2/2)$$

Thus for sufficiently large $m, v(t_m)$becomes negative which contradicts the Nonnegative nature of $v\ for\ t \geq 0$.

this contradiction shows that $|X_{1,}(\mathrm{t})| + |X_2(\mathrm{t})| \to 0$ as $\mathrm{t} \to \infty$ and the proof is complete.

As an example let us consider the case of two species Lotka- Volterra type competition in a periodic environment for for which we have

$$f_1\big(\mathrm{t}, x_1(\mathrm{t}), x_2(\mathrm{t} - T_{12})\big)1 \quad \left\{ \frac{x_1(\mathrm{t}) + ax_2(\mathrm{t} - T_{12})}{K_1(\mathrm{t})} \right\}$$

$$f_2(\mathrm{t}, x_1(\mathrm{t} - {}_{21})x_2(\mathrm{t})) = 1 - \left\{ \frac{\beta x_1(\mathrm{t}-T_{21}) + x_2(\mathrm{t})}{K_2(\mathrm{t})} \right\}$$

Where $\alpha, \beta$ are positive constants and $k_{1,}k_{2,}$ are strictly positive continuous periodic functions with a common period Equation (4.6)and (4.7) for this special case lead to

$$\left|\operatorname{axm}_{t>T}\{-a_{11}(t)\}\right| = \max_{t>T}(1/K_1(t)) > \min_{t>T}\left(-a_{21}(\mathrm{t})\right)\Big| = \min_{t\in}, \beta/K_2(\mathrm{t})$$

$$\max_{t\in T}\{-a_{22}(t)\}\Big| = \max_{t\in T}(1/K_2(t)) > \min_{t\in T}\left(-a_{12}(\mathrm{t})\right)\Big| = \min_{t\in} \mathfrak{a}/K_1(\mathrm{t})$$

The inequalities reduce in the absence of temporal variations in $k_{1,}k_{2,}$ to the inequality 0< $\mathfrak{a}$,$\beta$<1, which is well known for the global asymptotic stability of the positive steady state corresponding to the above example in a constant systemic domain the above example shows also that delays in interspecific interactions have no influence on the local asymptotic stability of the unique periodic solution of the periodic ecosystem A number of cases of this type in a constant environment have been established recently by the author (Gopalsamy [6]).A numerical example is illustrated at the end.

**Consequences**

It can be found that all the conclusions on the existence uniqueness and stability of a periodic solution of (2.1)-(2.4) are valid when $OT_{12}$=0 or $T_{12}$=0 or both $T_{12}$=0 and $T_{21}$=0: thus it will follow that delays in interspecific competition do not destabilize an otherwise stable coexistence (oscillatory or nonoscillatory) of two competing types of interacting particles atoms molecules transactions interaction. In forthcoming work we will extend the analysis of the system (2.1)-(2.4) to a system of n-competing species with a hypothesis similar to that of (2.4) It is however an open problem to investigate the effects of delays in intraspecific competition in a periodic milieu

Pianka ([8],pages 139-140) remarks that “In real population rate of increase competitive abilities and carrying capacities do very from individual to individual with population density in space and time Indeed temporal variation in the environment may often allow coexistence by continually altering the competitive abilities of population inhabiting it Moreover time lags are doubt lessly of some importance in real population Finally a heterogeneous environment may allow real competitors to evolve divergent resource utilization patterns and to reduce interspecific competitive inhibition

To understand the mechanism of selection pressure against competitive inhibition we first write the Malthusian parameters $\mu_1$, $\mu_2$ of fitness for the two species as follows (see MacArthur[7],page 226-230):

$$\mu_1=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_1\left(t,x_1(t),x_2(t-T_{12})\right)$$

$$\mu_2=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_2(t,x_1(t-T_{21}),x_2(t))$$

Since $\left(x_1(t),x_2(t)\right)\rightarrow\left(x_1^*(t),x_2^*(t)\right)$ for large t we can write the parameters of finess in the from

$$\mu_1=f_1\left(t,x_1^*(t),x_2^*(t-T_{12})\right)$$

$$\mu_2=f_2\left(t,x_1^*(t-T_{12}),x_2^*(t)\right)$$

In these relations we can use

$$x_2^*(t-T_{12}),x_2^*(t)-T_{12}\frac{dx_2^*(t-T_{12}+\theta_1T_{12}}{dt}$$

$$x_1^*(t-T_{12}),x_1^*(t)-T_{21}\frac{dx_1^*(t-T_{21}+\theta_2T_{21}}{dt}$$

$\theta_1,\theta_2,\in(0,1)$ and rewrite $\mu_1,\mu_2$ in the form

$$\mu_1=f_1\left(t,x_1^*(T),x_2^*(t)-T_{12}\frac{dx_2^*(t-T_{12}+\theta_1T_{12}}{dt}\right.$$

$$\mu_2=f_2\left(t,x_1^*(T),-T_{21}\frac{dx_1^*(t-T_{21}+\theta_2T_{21}}{dt},x_2^*(t)\right)$$

$\theta_1,\theta_2,\in(0,1)$

It is seen that the presence of the delays $T_{12}$ and $T_{21}$ provides additional sources of variability in the finess parameters and thereby in the selection pressure against interspecific competition Since a necessary and sufficient condition for natural selection to act is differential (variable) rates of confederation growth we find that time delays coupled with periodicity of an domain may allow real to evolve divergent resource utilization patterns and to reduce

Interspecific competitive inhibition for some remarks relating to the significance of a periodic environment to sustain the coexistence of a large number of confederation types in the absence of delays we refer to a recent article (Gopalsamy)The following example has been numerically solved and the solutions are displayed graphically for two different values of the delay parameter $\tau$:

$$\frac{dx_1(t)}{dt} x_1(t)[1 - x_1(t) + 0.4x_2(t - \mathcal{T})/k_1(t)]$$

$$\frac{dx_2(t)}{dt} x_2(t)[1 - \{0.8\, x_1(t - \mathcal{T}) + x_2(t)\}/k_2(t)]$$

Where $k_1(t) = 20 + \sin(\pi t)$ and $k_2(t)$=2.5+sin ($\pi$t)

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There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies

systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational

content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification.

## GOPALSAMY'S MODEL APPIED TO COSMOLOGICAL PROBLEMS

## GLOBAL ASYMPTOTIC STABILITY IN AN ALMOST PERIODIC LOTKA- VOLTERRA SYSTEM

### Abstract

**Gopalsamy studied** Sufficient conditions are obtained for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of a Lotka- Volterra system with almost periodic coefficients. Research was published Australian. Math Soc. B.27 (1986) 346-360. Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of

decay rate of perturbations. Sole intention, main objective and primary concern of the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are counted in the same vein as that of predator prey. Note that the model is applicable to each and every module. Quintessentially each parameter should have been given different name such as superscript 1,2, etc., as that has been done in the foregoing and a formal concatenation done at the end. Due to spatial restrictions we have not been able to do the line of thought and manner of analysis as has been done hereinbefore. This is the most important aspect of the model. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 1. Introduction

At the outset we assume that the system. This article is concerned with the derivation of a set of "easily verifiable" sufficient conditions for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of the Lotka- Volterra system

$$\frac{dx_1(t)}{dt} = x_i(t)\left\{b_1(t) - \sum_{j=1}^{n} a_{ij}\,(t)x_j(t)\right\};$$

$$t>t_0;\ t_0 \in (-\infty,\infty);\ i = 1,2,\dots,n, \qquad (1.1)$$

where $b_1, a_{ij}\ (i,j = 1,2,\dots.,n)$ are nonnegative almost periodic functions defined for $t \in (-\infty,\infty)$.

In question satisfies Haimovici's equation based on age wise stratification given in the foregoing. Note that Haimovici's equations can be applied to Lotka-Volterra system term by term.

Let R and $R_n$ respectively denote the set of all real numbers and the $n$ - dimensional real Euclidean space; $R_n^+$ will denote the nonnegative cone of $R_n$. Since almost- periodic functions defined on R are bounded, we define the constants $b_i^l, b_l^u, a_{ij}^l, a_{ij}^u, (i,j = 1,2,\dots,n)$ by the following:

$$\inf_{t\in R} b_l(t) = b_i^l; \qquad \sup_{t\in R} b_l(t) = b_l^u, \qquad i,j = 1,2,\dots,n.$$

$$\inf_{t\in R} a_{lj}(t) = a_{lj}^l; \qquad \sup_{t\in R} a_{lj}(t) = a_{ij}^u.$$

We shall study the almost- periodic system (1.1) with the following assumption on the coefficients of (1.1);

$$b_l^l > 0; \qquad a_{ll}^l > 0, \tag{1.2}$$

$$b_l^l > \sum_{\substack{j=1 \\ j\neq i}}^{n} a_{lJ}^u (b_j^u / a_{iJ}^l). \tag{1.3}$$

Since the solutions of (1.1) corresponding to nonnegative initial conditions remain nonnegative subsequently, it will follow that

$$\frac{dx_1}{dt} \leq x_l\{b_l^u - a_{ll}^l x_i\}; \qquad t > t_0, i = 1,2,\dots,n, \tag{1.4}$$

As a consequence of which we will have

$$0 < x_l(t_0) \leq b_l^u / a_{ll}^l = x_i^u \Rightarrow x_i(t) \leq b_l^u / a_{Il}^l, \qquad i = 1,2,\dots,n; t \geq t_0. \tag{1.5}$$

(1.1) And (1.5) together lead to

$$\frac{dx_1}{dt} \leq x_l\left\{b_l^u \sum_{\substack{j=1 \\ j\neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right) - a_{ll}^l x_i\right\}, \qquad i = 1,2,\dots,n; \tag{1.6}$$

Now (1.2), (1.3) and (1.6) lead to

$$x_i(t_0 \leq \left(b_i^l - \sum_{\substack{j=1 \\ J\neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right)\right)(1/a_{ll}^u) = x_i^l$$

$$\Rightarrow x_l(t) \to x_i^l \ for\ t \to t_0,\ i = 1, 2, \dots, n. \quad (1.7)$$

2. It is our aim to show that the mode together with other models stated in the foregoing together with haimovicis model is applicable to whole lot of consummative debit – credit systems in cosmology, astrophysics, finance, economics, sociology, anatomy computation. We will not however divide the system studied here into 9 components/ modules due to restrictions on space….

From the foregoing preparation we have the following:

LEMMA 1.1. *Assume that* $b_i, a_{ij} (i, j = 1, 2, \dots, n)$ *are scalar nonnegative almost periodic functions defined on* $\boldsymbol{R}$ *such that (1.2) and (1.3) hold. Then the set S defined by*

$$S = \{x = (x_1, x_2, \dots, x_n) \in \boldsymbol{R}_n | x_* \leq x_l \leq x^*; i = 1, 2, \dots, n\} \quad (1.8)$$

*Where*

$$x_* = \min_{1 \leq i \leq n} x_l^l; \qquad x_* = \min_{1 \leq i \leq n} x_i^u \quad (1.9)$$

*is invariant with respect to (1.1)*

**Appellation, cognomen, compellation, designation, denomination of the conducive, infelicitous and unfortutious** responses to environment change may be influenced by changes in available domain and milieu under consideration. An example is the following study: Authors describe the gravitational degrees of freedom of the Schwarzschild black hole by one free variable. Authors introduce an equation which we suggest to be the Schrödinger equation of the Schwarzschild black hole corresponding to this model. We solve the Schrödinger equation explicitly and obtain the mass spectrum of the black hole as such as it can be observed by an observer very far away and at rest relative to the black hole. Our equation implies that there is no singularity inside the Schwarzschild black hole, and that the black hole has a certain ground state in which its mass is non-zero. **Schrodinger equation of the Schwarzschild black hole Jarmo Mäkelä (Cambridge U E-Print: gr-qc/9602008 | PDF** Particles, axions, atoms, molecules, quarks, Majorana fermions, bosons et al in its confederation capacity and aggregation processes, type interactions and interactions between landscape dynamics. Following bears ample testimony, infallible observatory to the fact there might be changes in the quantum behaviour due to merger of black holes. Following example bears enough testimony, infallible observatory and impeccable demonstration and apodictic evidence as how some of the models could be applied to cosmological problems. This is attributable and ascribable to the fact that the classification criterion is based on the age.

**How gravity kills (e) Schrödinger's cat Theorists argue that warped space-time prevents (e) quantum superpositions of large-scale objects. Elizabeth Gibney 17 June 2015 Propositions are written in model able form and can easy be scheduled.**

(1) If the cat in Erwin Schrödinger's famous thought-experiment behaved according to quantum theory, it would be able to exist in multiple states at once: both dead and alive.

(2) Physicists' common explanation for why we don't (e) see such quantum superpositions — in cats or any other aspect of the everyday world — is interference from (e) the environment.

(3) As soon as a quantum object interacts with (e&eb) a stray particle or a passing field, it picks (e) just one state, collapsing (eb) into our classical, everyday view.

(4) But even if physicists could completely isolate (eb) a large object in (eb) a quantum superposition, according to researchers at the University of Vienna, it would still collapse into (e&eb) one state — on Earth's surface, at least.

(5) "Somewhere in interstellar space it could be that the cat has (e) a chance to preserve quantum coherence, but on Earth, or near any planet, there's (=) little hope of that," says Igor Pikovski. The reason, he asserts, is gravity.

(6) Pikovski and his colleagues' idea, laid out in a paper published in Nature Physics on 15 June1, is at present only a mathematical argument. But experimenters hope to test whether gravity really does collapse (eb) quantum superpositions, says Hendrik Ulbricht, an experimental physicist at the University of Southampton, UK. "This is a cool, new idea, and I'm up for trying to see it in experiments," he says. Assembling the technology to do so, however, may take as long as a decade, he says.

(7) Cinema-goers who saw the film Interstellar are already familiar with the basic principle behind the Vienna team's work. Einstein's theory of general relativity states that an extremely massive object causes clocks near it to (e) run more slowly because (e) its strong gravitational field stretches (e&eb) the fabric of space-time (which is why a character in the film aged only an hour near a black hole, while (e) seven years passed on Earth). On a subtler scale, a molecule placed nearer the Earth's surface experiences (e&eb) a slightly slower clock than one placed slightly further away.

(8) Because of gravity's effect on (e&eb) space-time, Pikovski's team realised that variance in a molecule's position will also influence (e&eb) its internal energy — the vibrations of particles within (eb) the molecule, which evolve (eb) over time.

(9) If a molecule were put in a quantum superposition of two places, the correlation between position and (e&eb) internal energy would soon cause (eb) the duality to 'decohere' to (e) the molecule taking just one path, they suggest.

(10) "In most situations decoherence is due to (e) something external; here it's as though the internal jiggling is interacting with (e&eb) the motion of the molecule itself," adds Pikovski.

**A practical limit**

(11) No one has yet seen this effect because other sources of decoherence — such as magnetic fields, thermal radiation and vibrations — are (=) typically much stronger and cause (eb) quantum systems to collapse long before gravity becomes an issue. But experimenters are keen to try.

(12) Markus Arndt, an experimental physicist also at the University of Vienna, already tests whether quantum superpositions can be observed for (e) large, though not cat-sized, objects (see Fattening up Schrödinger's cat).

(13) He sends large molecules through (e&eb) a matter-wave interferometer, a system that gives (eb) each molecule a choice of (e) two different paths.

(14) In the classical view, a molecule travels down just one path; a quantum molecule effectively passes through (e&eb) both routes at once and interferes with (e&eb) itself to create (eb) a characteristic wave-like pattern (see image, below).

| |
|---|
| (15) A quantum interference pattern built up from complex molecules known as phthalocyanines.(deleted due to spatial constraints) |
| (16) A similar set up could be used to (e) test for gravity's ability to destroy (e) quantum behaviour: by comparing a vertical interferometer, in which the superposition should soon decohere due to (e) the stretching of time in one path versus (e&eb) the other, with a horizontal set up, where (e) the superposition could remain. |
| (17) Arndt, who has tested the effect for (e) molecules of up to 810 atoms2, points out those large molecules, would be (=) good for testing the gravitational effect because (e) they contain (e) lots of particles contributing to (e) internal energy. |
| (18) **But not only would the researchers have to further suppress the external environment to reduce (e) other decoherence effects, they would also need to (e) increase the separation of the two paths from micrometres to metres, or else use molecules perhaps 1 million times more massive. "It's certainly very challenging," says Arndt**. |
| (19) If gravity's effect does limit (e) quantum behaviour on Earth, tests of quantum reality for (e) large objects may eventually have to (e) move into space, says Angelo Bassi, a physicist at the University of Trieste in Italy. |
| (20) "But from a deep, fundamental point of view, this is nothing new," he says. A gravitational field is merely another environment to interact with, so invoking it does not (e) explain whether quantum behaviour might lead to (eb) classical reality if gravity's influence were mitigated — for example, by doing the experiment in (eb) gravity-free space. |
| (21) Nor does the effect described by Pikovski and colleagues say anything about quantum gravity: a theory that would unite gravity and (e&eb) quantum mechanics in (eb) a single description, which many researchers are working towards. |
| (22) "It's an interesting effect, but it's still quantum physics applied to (e&eb) classical general relativity. In that way it doesn't change our picture of the world," adds Bassi. Nature doi:10.1038/nature.2015.17773 |

Integrated factors these influences because they deal with ambience, mise en scene change or confederation dynamics, but rarely both. Time series of (e) space time suitability models with spatially explicit stochastic models to explore factors that influence the viability of type viability and confederation feasibility under stable and changing scenarios. Methodology could be applied to interrogation propositions and problems which might subsume one with the other notwithstanding that there is a space of singularities and there is space for completion of present from the past and future. An example is Time image movement from Deleuze. We clearly face a stiff challenge in our attempt to proffer Deleuzian treatments of photography as our critical project brushes up against the fact that Deleuze does not invest photography with (e&eb)the creative and re-creative potential that he associates with (e&eb) both cinema and modern painting. In short, while he highly values other visual artistic forms, he seemingly presents photographic texts as stagnate documents or tools that produce certainty, organize bodies and desires, and iterate hackneyed ideas. He even uses photography as something of a foil to demonstrate the innovation of the cinema and the originality of modern painters. In Francis Bacon: The Logic of Sensation (1981), he notes how "photography has taken over the illustrative and documentary role, so that modern painting no longer needs to fulfill this function, which still burdened earlier painters" (10). He later explains that "photographs are ways of seeing, and as such, they are illustrative and narrative reproductions or representations ... . They are what are seen, until finally one sees nothing else." He treats the photo as an instrument for reproducing representations of reality—a device that iterates images until they are ossified as established stories, icons, or even stagnant perceptions. Indeed, like Bergson before him, Deleuze compares photography to what he identifies as a reductionist model of perception: "what we see, what we perceive, are photographs" (74).[4] The photo is, in effect, an always-already passé sensational experience that numbs our sensitivity to the ongoing vitality of life and provides us with an efficient

understanding of the dynamism of our world. But it is compelling that Deleuze theorizes both the quotidian nature and the liberatory potential of photography; he treats it as a common, everyday visual experience that, despite its aesthetic and technological limitations, emancipates the modern painter to explore new artistic opportunities. This special issue of Rhizomes repeatedly provides artistic examples and critical strategies that invite us to see photography as replete with the power to unburden and emancipate that Deleuze associates with both film and painting. Indeed, one way to think of this project of reading Deleuze against Deleuze might be to ask: if photography can free painting and painters to pursue novel creative opportunities, can it likewise free itself, photographic practitioners, and even viewers to create and re-create anew? **Rhizomes » Issue 23 (2012) » Michael Kramp Unburdening Life, or the Deleuzian Potential of Photography Michael Kramp Lehigh University** dmk209@lehigh.edu

**Existence of an almost- periodic solution**

DEFINITION 2.1. A function $f(t,x)$,where $f$ is an $m$- vector, $t$ is a real scalar and $x$ is an $n$- vector, is said to be almost periodic in $t$ uniformly with respect to $x \in X \subset R_n$, $if\ f(t,x)$ is continuous in $t \in R\ and\ x \in X$, and if for any $\varepsilon > 0$, it is possible to find a constant $l(\varepsilon) > 0$ such that in any interval of length $l(\varepsilon)$ there exists a $\tau$ such that the inequality

$$\left|\left|f\ (t+\ \tau,x) - \ f(t,x)\right|\right| = \sum_{l=1}^{m} |f_l\ (t+\ \tau,x) = \ f_l(t,x)|$$

$$< \varepsilon \qquad (2.1)$$

is satisfied for all $t \in (-\infty,\infty), x \in X$. The number $\tau$ is called an $\varepsilon -$ translation number of $f(t,x)$.

Let $\{\lambda_j\}$ denote the set of all real numbers such that

$$\lim_{T\to\infty} T^{-1} \int_0^T f(t,x)\exp(-i\,\lambda_j\,t)dt \not\equiv 0 \quad for\ x \in X; i = \sqrt{-1}. \qquad (2.2)$$

It is known that when the set X is separable, the set of numbers $\{\lambda_j\}$ in (2.2) is countable. The set $\{\sum_1^N n_j\ \lambda_j\}$ for all integers N and integers $n_j$ is called the module of $f(t,x)$.

The following result concerned with module containment will be used in the proof of our existence theorem below.

**MODULE CONTAINMENT THEOREM** (Yoshizawa [7], page 18). *Let $f(t,x)$ and $g(t,x)$ be almost periodic in t uniformly for $x \in D \subset \boldsymbol{R}_n$. If for any compact set $S \subset D$ and for any sequence of real numbers $\{\tau_k\}$ having a limit as $k \to \infty$ (the limit being finite or$\pm\infty$) for which $\{f(t+\tau_k,x)\}$ is uniformly convergent on $\boldsymbol{R} \times$*

$S, \{g(t+,x)\}$ *is also uniformly convergent on* $\boldsymbol{R} \times S$ *then the module of* $g(t+,x)$ *is contained in the module of* $f(t+,x)$.

We can now formulate our existence theorem.

**THEORM 2.1** *Assume that the almost-periodic parameters* $b_1, a_{1J}$ $(i,j = 1,2,\dots,n)$ *of (1.1) satisfy (1.2)-(1.3). Furthermore assume that there exists a positive number* $\mu$ *such that*

$$\inf_{1 \in \boldsymbol{R}} a_{11}(t) > \sum_{\substack{J=1 \\ J \neq 1}}^{n} (\sup_{t \in \boldsymbol{R}} a_{Jl}(t)) + \mu, \quad i = 1,2,\dots,n. \tag{2.3}$$

*Then the system (1.1) has a strictly positive (component wise) almost-periodic solution, say* $w(t) = w_1(t), \dots, w_n(t), t \in \boldsymbol{R}$, *whose module is contained in that of*

$$f(t,x) = \{f_1(t,x_1,\dots,x_n), f_2(t,x_1,\dots,x_n), \dots, f_n(t,x_1,\dots,x_n)\}$$

*Where*

$$f_l(t,x_1,\dots,x_n) = x_l \left\{ b_l(t) - \sum_{j=1}^{n} a_{lj}(t) x_J \right\},$$

$$i = 1,2,\dots n; (x_1,x_2,\dots,x_n) \in S. \tag{2.4}$$

PROOF. The set $S$ defined by (1.8)-(1.9) is compact in $\boldsymbol{R}_n$. Let $\{\tau_h\}$ be any sequence of real numbers such that $\tau_k \to \infty$ as $k \to \infty$ and

$$\sum_{l=1}^{n} |f_i(t+\tau_k, x_1, \dots, x_n) - f_l(t, x_1, \dots, x_n)| \to 0 \quad \text{as } k \to \infty \tag{2.5}$$

Uniformly in $(x_1, \dots, x_n) \in S$ for $t \in (-\infty, \infty)$. for any real number $\beta$let $k_0(\beta)$ be the smallest value of $k$ such that $\tau k_0 + \beta \geq t_0$. Since the invariance of $S$ implies that any solution of (1.1)

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), x_2(t+\tau_k, t_0, x_2(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\} \in S$$

For all $t \geq \beta$ and $k \geq k_0(\beta)$ whenever $\{x_1(t_0), \dots, n_n(t_0)\} \in S$, our first task in the proof is to show that the sequence

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Converges to a continuous bounded function say $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined for $t \in [\beta, \infty)$ such that $w_1(t) \in [x_*, x^*], t \in [\beta, \infty), i = 1, 2, \dots, n$ and the convergence is uniform on all compact subsets of $[\beta, \infty)$. Since $x_*$ and $x^*$ are independence of $\beta$, it is enough to show that

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Is a Cauchy sequence on compact subsets of $[\beta, \infty)$.

Let $U$ be any compact subset of $\beta, \infty)$., and let be any arbitrary positive number. Choose an integer $n_0(\varepsilon, \beta) \geq k_0(\beta)$ so large that for $m \geq k \geq n_0(\varepsilon, \beta)$

and all $t \in (-\infty, \infty)$

$$\sum_{i=1}^{n} |b_1(t+\tau_m) - b_1(t+\tau_k)| < (\mu\varepsilon x_*)/(4x^*), \tag{2.6}$$

$$x^* \sum_{i=1}^{n}\sum_{j=1}^{n} |a_{lj}(t+\tau_m) - a_{lj}(t+\tau_k)| < (\mu\varepsilon x_*)/(4x^*) \tag{2.7}$$

$$(2nx^*/x_*)\exp[-\mu x_*(\beta + \tau_k - t_0)] < (\varepsilon/2x^*)$$

(2.8)

Denote $x_i(t, t_0, x_1, (t_0))$ $by$ $x_1(t)$ brevity, and consider a function $v(s)$ defined by

$$v(s) = v\left(x(s), x(s+\tau_m+\tau_k)\right)$$

$$= \sum_{i=1}^{n} |\log x_1, (s) - \log x_1(s+\tau_m+\tau_k)|, \qquad x \geq t_0, \tag{2.9}$$

Where $x(t) = \{x_1(t), \dots, x_n(t)\}$ denotes a solution of (1.1) with $\{x_1, (t_0), \dots, x_n(t)\} \in S$. the invariance of S for (1.1), together with the elementary mean value theorem of differential calculus, implies

$$(1/x^*)\sum_{l=1}^{n}|x_l(s)-x_1(s+\tau_m+\tau_k)|\ \upsilon(s)$$

$$(1/x^*)\sum_{i=1}^{n}|x_l(s)-x_1(s+\tau_m+\tau_k)|. \tag{2.10}$$

Calculating the right derivative $D^+\upsilon$ of $\upsilon$ along the solutions of (1.1) and simplifying

$$D^+\upsilon(s)=\sum_{i=1}^{n}\{sign[x_i(s)-x_i(s+\tau_m-\tau_k)]\}$$

$$\left\{\frac{d[logx_i(s)]}{ds}-\frac{d[logx_i(s+\tau_m-\tau_k)]}{ds}\right\}$$

$$\sum_{l=1}^{n}|b_l(s)-b_l(s+\tau_m-\tau_k)|$$

$$+\sum_{l=1}^{n}(-a_{il}(s)|x_i(s)-x_i(s+\tau_m-\tau_k)|+\sum_{\substack{J=1\\ j\neq l}}^{n}a_{ij}(s)|x_j(s)-x_j(s+\tau_m-\tau_k)|)$$

$$+\sum_{i=1}^{n}\sum_{j=1}^{n}|a_{ij}(s)|a_{ij}(s+\tau_m-\tau_k)|\,x_j(s+\tau_m-\tau_k) \tag{2.11}$$

$$-\mu x_*\upsilon\big(x(s),x(s+\tau_m-\tau_k)\big)+\sum_{l=1}^{n}|b_l(s)-b_i(s+\tau_m-\tau_k)|$$

$$+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|. \tag{2.12}$$

Let us set $s-\tau_k=\sigma$; then we have from (2.6)-(2.7) that

$$\sum_{i=1}^{n}|b_l(s)-b_1(s+\tau_m-\tau_k)|+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|.$$

$$= \sum_{l=1}^{n} |b_l(\sigma + \tau_k) - b_i(\sigma + \tau_m)|$$

$$+x^* \sum_{i=1}^{n} \sum_{J=1}^{n} |a_{ij}(\sigma + \tau_k) - a_{ij}(\sigma + \tau_m)|$$

$$< (\mu\varepsilon x^*)/(2x^*). \quad (2.13)$$

Now an integration of (2.12) over $(t_0, t - \tau_k)$ together with (2.13) leads to

$$v\big(x(t+\tau_k), x(t+\tau_m)\big) \leq v\big(x(t_0), x(t_0 + \tau_m - \tau_k)\big) \exp[-\mu x_*(t + t_k - t_0)]$$

$$+[(\mu\varepsilon x_*)/(2x^*)]/\mu x_*)$$

$$\leq [(2x_n^*)/x_*]\exp[-\mu x_*(t + t_k - t_0)] + (\varepsilon/2x^*)$$

$$< (\varepsilon/x^*) \text{ On using (2.8), (2.10)} \quad (2.14)$$

For $m \geq n_0(\varepsilon, \beta)$ and for all $t \in \cup \subset [\beta, \infty]$, Thus we have from the definition of $v$ and (2.14) that

$$\sum_{l=1}^{n} |log x_i(t + \tau_k) - log x_l(t + \tau_m)| < (\varepsilon/x^*) \quad (2.15)$$

And hence using $x_* \leq x_i(t) \leq x^*, i = 1, 2, \dots, n; t \geq t_0$,

$$\sum_{i=1}^{n} |x_i(t + \tau_k) - x_i(t + \tau_m)| < \varepsilon \text{ for all } t \in \cup \subset [\beta, \infty], m \geq n_0(\varepsilon, \beta). \quad (2.16)$$

The existence of a function $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined on $[\beta, \infty]$, such that $w_i(t) \in [x_*, x^*], i = 1, 2, \dots, n; t \in [\beta, \infty]$ follows from (2.16) and the fact that $x_i(t) \in [x_*, x^*], i = 1, 2, \dots, n;$ for $t \geq t_0$. Since $\beta$ is arbitrary, $w$ is in fact

Defined on $(-\infty, \infty)$ and hence we have

$x_i(t+\tau_k) \to w_l(t)$ as $k \to \infty$ uniformly on compact subsets of $(-\infty,\infty)$, $i = 1,2,\dots,n.$

(2.17)

Our next task is to show that $w: R \to S$ is differentiable and is a solution of (1.1). Since $x(t) = \{x_1(t),\dots,x_n(t)\}$ satisfies (1.1) we have $\sum_{i=1}^{n} |\frac{d}{dt} x_i(t+\tau_k) - \frac{d}{dt} x_i(t+\tau_m)|$

$$\sum_{i=1}^{n} |f_l\left(t+\tau_k, x_1(t+\tau_k),\dots,x_n\left(t+\tau_k\right)\right)$$

$$- f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)| \qquad (2.18)$$

$$+\sum_{l=1}^{n} |f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$- f_l\left(t+\tau_m, x_1(t+\tau_m),\dots,x_n(t+\tau_m)\right)|$$

Where $f_l\ (i = 1,2,\dots,n)$are as in (2.4). Since $x(t+\tau_k) = x_1(t+\tau_k),\dots, x_n(t+\tau_k) \in S$ for large $\tau_k$ for $t$ in each compact subset of $(-\infty,\infty)$ there exists an integer $n_1(\varepsilon) > 0$ such that if $m \text{ } k \text{ } n_1(\varepsilon)$, then for any $\varepsilon > 0$,

$$\sum_{i=1}^{n} |f_l\left(t+\tau_k, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$-f_1\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)| < (\varepsilon/2). \qquad (2.19)$$

Similarly since $x_l(t+\tau_k) \in S$ and $x_l(t+\tau_m) \in S$ for $i = 1,2,\dots,n$ there is an integer $n_2(\varepsilon) > 0$ such that for large enough $m, k$, such that $m \text{ } k \text{ } n_2(\varepsilon)$, we have

$$\sum_{l=1}^{n} |f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$-f_1\left(t+\tau_m, x(t+\tau_m),\dots,x_n(t+\tau_m)\right)| < (\varepsilon/2) \qquad (2.20)$$

This equation will be very much useful in chemistry if used with the modules mentioned and written in concatenation with each other. It should be borne in mind that the applicability is modulewise and in no case concatenation should be marginalized. In all cases parameters are different which could be represented with different superscripts and subscripts like 1,23,…….etc., **Anal Bioanal Chem. 2004 Sep; 380(1):36-45. Epub 2004 Jul 29. Electronic tongue for pharmaceutical analytics: quantification of tastes and masking effects Legin A et al**

The organoleptic aspects of pharmaceutical formulations affect their acceptability to the patient and hence can have an important effect on concordance with treatment. Objective evaluation of these aspects, particularly the taste of the formulation and the drug substance it contains, is difficult. Whilst volunteer taste panels can be used to good effect their utility is limited, particularly during very early stage development when the toxicological profile of the active pharmaceutical ingredient (API) is yet to be established in detail. A potentiometric "electronic tongue" has been applied to analyse a variety of 41 individual substances and mixtures of particular interest for pharmaceutical research and development. The electronic tongue (ET) was capable of discriminating between substances with different taste modalities and could also distinguish different substances eliciting the same basic taste; the ET is promising in terms of quantifying the content of each substance and has an ability to detect nuances of the basic taste (e.g. lingering or short-lived). After calibration the electronic tongue was successfully applied to predicting bitterness strength of binary mixtures with a sweetener in terms of "apparent" or "perceived" quinine content. In order to render a formulation palatable it is often necessary to mask the (usually bitter) taste of the API by the addition of masking agents such as sweeteners and flavours. The ET proved capable of distinguishing between formulations with different levels of sweetener and/or flavour in a manner that was consistent with their masking efficiency as perceived by a small human taste panel. A suitably calibrated ET could have the benefit of providing the pharmaceutical formulator with reliable data concerning the taste of the product quickly and with a reduced need to ask volunteers to taste active pharmaceutical samples. Early development activities could be facilitated when human tasting is usually not possible in the absence of the required toxicological data.

Thus if $m \; k \; n_0(\varepsilon) = max\{n_1, n_2\}$, we have

$$\sum_{l=1}^{n} |\frac{d}{dt}\, x_1(t+\ \tau_k) - \frac{d}{dt} x_1\ (\ (t+\tau_m)| < \varepsilon \qquad (2.21)$$

Which shows that

$$\lim_{k\to\infty} \frac{d}{dt} x_1(t+\ \tau_k) \quad exists\ for\ i = 1, 2, \dots, n$$

Uniformly on all compact subsets of $(-\infty, \infty)$. Thus we have

$$\lim_{k\to\infty} \lim_{k\to\infty} \frac{x_i(t+\tau_k+h) - \ x_i(t+\tau_k)}{h}$$

$$= \lim_{h\to 0} \lim_{k\to 0} \frac{x_l(t+\tau_k+h) - \ x_l(t+\tau_k)}{h} \qquad (2.22)$$

$$= \lim_{h\to 0} \frac{w_1(t+h) - \ w_1(t)}{h}, \qquad i = 1, 2, \dots, n,$$

Showing that (since the left side of (2.22) exists) $w_i (i = 1, 2, \dots, n)$ is differentiable. Furthermore we have from (2.22) that

$$\frac{dw_l(t)}{dt} = \lim_{k\to\infty} [\, f_i\ \big(t+\tau_k, x_1(t+\tau_k), \dots, x_n\ (t+\tau_k)\big)$$

$$-f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big)$$

$$+ - f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big)]$$

$$= \ -f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big); \qquad i = 1, 2, \dots, n, \qquad (2.23)$$

By the uniform almost periodicity of $f_1, f_2, \dots, f_n$ on $R \times S$ and hence $w(t) = \big(w_1(t), \dots, w_n(t)\big)$is a solution of (1.1).

In order to complete the proof, we have to show that $w(t) = \big(w_1(t), \dots, w_n(t)\big)$ is an almost- periodic function whose module is contained in that of $\{f_1(t, x_1, \dots, x_n), f_2(t, x_1, \dots, x_n), \dots, f_n(t, x_1, \dots, x_n)\}$ for $(x_1, x_2 \dots, x_n) \in S$. For this purpose it is in fact sufficient (by the module containment theorem) to show that for any sequence of real numbers $\{\tau_k\}$ for which

$$\{f_1(t + \tau_k, x_1, \dots, x_n), f_2(t, + \tau_k, x_1, \dots, x_n), \dots, f_n(t, + \tau_k, x_1, \dots, x_n)\}$$

Converges uniformly on R for all $(x_1, \dots, x_n) \in S$, the sequence $\{w_1(t, + \tau_k) \dots, w_n(t + \tau_k)\}$ converges uniformly on R where $\tau_k$ tends to a finite number or $\pm\infty$ as $k \to \infty$.

In the case where $\tau_k \to \tau \neq \pm\infty$ as $k \to \infty$, we have $w_l(t, + \tau_k) \to w_i(t + \tau), i = 1, 2, \dots, n$ uniformly on R as $k \to \infty$. Let us suppose $\tau_k \to \infty$ as $k \to \infty$. For any $\varepsilon > 0$ there exists an integer $n_0(\varepsilon)$ such that if $m$ [illegible] $n_0(\varepsilon)$ then

$$\sum_{l=1}^{n} |b_l(t + \tau_k) - b_l(t + \tau_m)| < (\varepsilon/4x^*)x_*, \qquad (2.24)$$

$$x^* \sum_{i=1}^{n} \sum_{j=1}^{n} |a_{ij}(t + \tau_k) - a_{ij}(t + \tau_m)| < \left(\frac{\varepsilon}{4x^*}\right)\mu x_* \qquad (2.25)$$

$$n(2x^*/x_*) \exp[\mu x_* \tau_k] < (\varepsilon/2x^*). \qquad (2.26)$$

Consider the function $v\big(w(s), w(s + \tau_m - \tau_k)\big)$ for $s \in [t, t + \tau_k]$ defined by

$$v\big(w(s), w(s + \tau_m - \tau_k)\big) = \sum_{l=1}^{n} |\log w_l(s) - \log w_l(s + \tau_m - \tau_k)|. \qquad (2.27)$$

Using the fact $w: \boldsymbol{R} \to S$ and $w$ is a solution of (1.1), calculate the right derivative $D^+v \ of \ v$;

$$D^+v\big(w(s), w(s + \tau_m - \tau_k)\big) = \sum_{i=1}^{n} \{\text{sign}[w_l(s) - w_i(s + \tau_m - \tau_k)]\}.$$

$$\left\{\frac{dw_i(s)}{ds} \cdot \frac{1}{w_l(s)} - \frac{dw_l(s + \tau_m - \tau_k)}{ds} \frac{1}{w_l(s + \tau_m - \tau_k)}\right\}$$

$$\leq -\mu \sum_{l=1}^{n} |w_l(s) - w_l(s+\tau_m-\tau_k)| + \sum_{l=1}^{n} |b_i(s) - b_l(s+\tau_m-\tau_k)|$$

$$+x^* \sum_{l=1}^{n}\sum_{J=1}^{n} |a_{lJ}(s) - a_{lJ}(s+\tau_m-\tau_k)| \quad (2.28)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+ \sum_{i=1}^{n} |b_l(s) - b_l(s+\tau_m-\tau_k)| \quad (2.29)$$

$$+x^* \sum_{l=1}^{n}\sum_{j=1}^{n} |a_{lJ}(s) - a_{ij}(s+\tau_m-\tau_k)|.$$

If we set $s-\tau_k = \sigma$ then we have from (2.29),

$$D^+ v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+ \sum_{l=1}^{n} |b_l(\sigma+\tau_k) - b_l(s+\tau_m)|$$

$$+x^* \sum_{l=1}^{n}\sum_{j=1}^{n} |a_{ij}(\sigma+\tau_k) - a_{ij}(\sigma+\tau_m)| \quad (2.30)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+[(\varepsilon\mu x_*)/(4x^*)] + [(\varepsilon\mu x_*)/(4x^*)]$$

(2.31)

Integrating (2.31) over $(t, t+T_k)$

$$v(w(t+T_k)w(t+T_m))$$

Another MODEL where the model is directly applicable with the scheduled ones are the case of insurance and corporate finance. Insurance: Mathematics and Economics Volume 42, Issue 1, February 2008, Retrieval of Black–Scholes and generalized Erlang models by perturbed observations at a fixed time Dedicated to Professor André Dubey on the occasion of his 60th birthday Daniel Neuenschwander

S-stable laws on the real line (more generally on Hilbert spaces), associated with some non-linear transformations (so-called "shrinking operations"), were introduced in [Jurek, Z.J., 1977. Limit distributions for truncated random variables. In: Proc. 2nd Vilnius Conference on Probability and Statistics, June 28–July 3, 1977. In: Abstracts of Communications, vol. 3, pp. 95–96; Jurek, Z.J., 1979. Properties of s-stable distribution functions. Bull. Acad. Polon. Sci. Sér. Math. XXVII (1), 135–141; Jurek, Z.J., 1981. Limit distributions for sums of shrunken random variables. Dissertationes Math. vol. CLXXXV]. In [Jurek, Z.J., Neuenschwander, D., 1999. S-stable laws in insurance and finance and generalization to nilpotent Lie groups. J. Theoret. Probab. 12 (4), 1089–1107], the authors interpreted s-stable motions on the real line as limits of total amount of claims processes (up to a deterministic premium) of a portfolio of excess-of-loss reinsurance contracts and showed that they led to Erlang's model or to Brownian motion. In [Neuenschwander, D., 2000b. On option pricing in models driven by iterated integrals of Brownian motion. In: Mitt. SAV 2000, Heft 1, pp. 35–39], we considered stochastic integrals whose integrand and integrator are both independent Brownian motions, thus modelling a stochastic volatility; as a result we got an analogue of the Black–Scholes formula in this model, confirming a result of Hull and White [Hull, J., White, A., 1987. The pricing of options on assets with stochastic volatility. J. Finance XLII (2), 281–300]. In the present paper, we will look at a common generalization of these processes, namely s-stable motions on the real line perturbed by a stochastic integral whose integrand and integrator are both (not necessarily independent) s-stable motions. The main result will be that if we can observe the distribution of such so-perturbed s-stable motions (together with the values of the perturbing processes) at time t=1, then we can identify the whole model (including the perturbation) among all models with Lévy processes perturbed by an iterated stochastic integral of two Lévy processes (in the Gaussian case) resp. among all models with a compound Poisson process with drift perturbed by an iterated stochastic integral of two compound Poisson processes (in the completely non-Gaussian case if the perturbing processes have no drift) without knowing anything about the history or about its distribution during $0 \le t < 1$. This applies, e.g., to a situation

where several assets obey the same model and one can estimate the distribution at time one by looking at the values of all these assets at time t=1. Interestingly enough, it will be convenient to treat the whole matter in the algebraic framework of the so-called Heisenberg group. This is a concept coming in fact from quantum mechanics and is in a certain sense the simplest non-commutative Lie group.

$\leq e^{-\mu x^* T_k} v\big(w(t), w(t+T_m-T_k)\big) + (\varepsilon/2x^*)$

$\leq e^{-\mu x^* T_k} \sum_{i=1}^{n} |\log\ w_i(t) - \log\ w_i(t+T_m-T_k)| + (1/x_*) + (\varepsilon/2x^*)$

$\leq n(2x^*/x_*)\exp[-\mu x_* T_k] + (\varepsilon/2x^*)$

$<(\varepsilon/x^*)$.

Since

$$v\big(w(t+T_k)w(t+T_m)\big) \leq \sum_{i=1}^{n} |\ w_1(t+T_k) - w_1(t+T_m))/x^*)$$

We have from(2.33)that

$$\sum_{i=1}^{n} |\ w_1(t+T_k) - w_1(t+T_m)| < \varepsilon\ for\ m \geq k \geq n_0(\varepsilon)$$

Which show that the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\}$converges uniformly on **R**. By a similar argument the above conclusion is also valid when $T_k \to -\infty\ as\ k \to \infty$. Thus the uniform convergence of the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\}\ as\ k \to \infty\ on\ \mathbf{R}$ follows and this completes the proof of the theorem.

**COROLLARY.** Suppose $b_i a_{ij}$ (i, j = 1.2, … , n)in (1.1)are periodic with a common period say $T > 0$. *Then if* (1.2), (1.3)*and* (2.3)*hold the system* (1.1)*will have a periodic solution* $w(t) = \{w,(t), \ldots\ldots, w_n(t)\}$*of period* $T$ *such that*

$$w_1(t+T) = w_1(t), \qquad t \in \boldsymbol{R}: i = 1,2, \ldots, n.$$

**PROOF.** Proof follows from that of Theorem 2.1 if we Choose $T_k = KT(K = 0,1,2, \ldots)$*and note that*

$$x_1(t+KT) - w_1(t) \to 0\ as\ k \to \infty; i = 1,2, \ldots, n; t \in \boldsymbol{R},$$

$$x_i(t+(K+1)t) - w_1(t) \to 0 \ \ as \ \boldsymbol{K} \to \infty; i = 1,2,\dots,n; t \in \boldsymbol{R},$$

$$x_i(t+T+KT) - w_1(t+T) \to 0 \ \ as \ k \to \infty; i = 1,2,\dots,n; t \in \boldsymbol{R},$$

Implying that

$$w_1(t+T) = w_i(t); \quad i = 1,2,\dots,n; t \in \boldsymbol{R},$$

## 3. Global asymptotic stability

By definition we say that an almost-periodic solution say $w(t) = \{w_1(t),\dots,w_n(t)\} of (1.1) is\ globally\ asymptotically$ stable (or attractive) if and only if every other solution x(t)={ $x_1(t),\dots,x_n(t)$}with $x_1(o) > 0, i = 1.2,\dots,n$ is defined for all t□0 and satisfies

$$\lim_{t\to\infty} |w_1(t) - x_1(t)| = 0; \quad i = 1,2,\dots,n.$$

A consequence of such a global asymptotic stability is that there cannot be another strictly positive almost-periodic solution of (1.1)

**THEOREM 3.1** Suppose the almost- Periodic coefficients $b_1$, $a_{ij}$ (i, j = 1,2, … , n)

Satisfy (1.2)and (1.3)and the following

$$\inf_{t\in R} a_{ij}(t) > \left\{ \sum_{\substack{J=1 \\ J\neq 1}}^{n} \sup a_{ij}(t) \right\} + \mu; \qquad i = 1,2,\dots,n$$

For some positive number $\mu\ Then$ (1.1)has an almost-periodic solution (with strictly positive components) which is globally asymptotically stable.

We will omit the proof since the proof is identical with that of Theorem 3.1 of [3]

Under condition (3.2) the almost –periodic solution $w(t) = \{w_1(t), \ldots . w_n(t)\}$ of (1.1)enjoys a somewhat strong from of stability even if the coefficients of (1.1)are subjected to a class of perturbations We note (see Yoshizawa [7]page17)that if a function b:R→ $R$ is almost periodic then H(b),the "hull of b" is defined as follows:

H(b)= g|g:**R**→ $\boldsymbol{R}$ $and\ g(t)\lim_{k\to\infty} b(t+T_k) for$ some sequence

$T_k$ of real numbers the limit being uniform in $t \in R$.

By Theorem 2.2 of Yoshizawa [7, page 10], members of the hull of almost periodic functions are themselves almost periodic.

A system of the from

$$\frac{dy_1(t)}{dt} = y_1(t)\left\{B_i(t) - \sum_{J-1}^{n} A_{ij}\ (t)y_1(t)\right\},$$

*i=1,2,…,n; t>*$t_0 \in R$,

Where

$$B_i \in H(b_i); (A_{ij}), \qquad \text{i, j} = 1,2, \ldots, \text{n},$$

is known as a system from the hull of (1.1) As a consequence of the definition of the hull in (3.3)and (3.5)we shall have

$$\inf_{t\in R} B_i(t) = \text{b}_i^l; \quad \sup_{t\in R} B_i(t) = \text{b}_i^l;$$

$$\inf_{t\in R} A_{ij}(t) = \text{b}_i^l; \quad \sup_{t\in R} A_{ij}(t) = \text{a}_{ij}^u;$$

and hence the coefficients of(3.4)- (3.5)satisfy the same conditions as (1.2),(1.3)and (3.2)It will then follow from theorems(2.1)and (3.1)that solutions of (3.4)-(3.5)with $y_1(t_0) > 0, i = 1,2, \ldots, n\ will\ have$ the following property:

$$y_1(t) > 0 \ \ for\ t \geq t_0; i = 1,2,\dots,n,$$

$$\sum_{i=1}^{n} y_1(t) remains\ bounded\ for\ t \geq t_0$$

**DEFINITION** A Solution $x(t) = \{x_1(t), \dots, x_n(t)\}, t \in t_0 \in \boldsymbol{R}$,of (1.1)is said to be stable under perturbations from the hull of (1.1)if the following holds :let $y(t) = \{(y_1(t), \dots, y_n(t)\}$ be any solution of (3.4)-(3.5)such that $y_1(t_0) > 0$, i=1,2,…,n,and for some positive constants $\beta_1, \beta_2 \beta_1, \leq y_1(t) \leq \beta_2, t \geq t_0, \ \ i = 1,2,\dots,n$ Then given any $\varepsilon > 0$,there exists a constant $\delta > 0$ such that

$$\sum_{j=1}^{n} |x_1(t_0) - y_1(t_0)| < \delta$$

$$\sum_{i=1}^{n} \sup_{t \in R} |b_1(1) - B_1(t)| < \delta \quad = \sum_{i=1}^{n} |x_1(t) - y_i(t)| < \varepsilon$$

$$\sum_{i=1}^{n} \sum_{j=1}^{n} | \quad \sup_{t \in R} |a_{ij}\ (t) - A_{ij}(t)| < \delta$$

For $t \geq t_0$.

We can now formulate the following result which shows that all solution of (1.1)are stable under perturbations from the hull of (1.1)

**THEOREM 3.2**. Assume that the conditions of Theorem (3.1) hold. Then every strictly positive solution of (1.1) ( including its unique almost-periodic solution) is stable under perturbations from the hull of (1.1).

PROOF. Let $x(t) = \{x_1(t), \dots, x_n(t)\}$ and $y(t) = \{y_1(t), \dots, y_n(t)\}$ be any two strictly positive solutions of (1.1) and (3.4) – (3.5) respectively for $t \geq t_0 \in R$. Let $\beta_1, \beta_2$ be constants such that

$\beta_1 \leq x_i(t) \leq \beta_2$ ,
$\beta_1 \leq y_i(t) \leq \beta_2$ , $i = 1,2,\dots,n; t \geq t_0$ (3.7)

Consider now a function $V(t) = V(x(t)y(t))$ defined by

$$V(x(t)y(t)) = \sum_{i=1}^{n} |\log x_1(t) - \log, y_i(t)|, \quad t \geq t_0 \in \text{R}. \tag{3.8}$$

It is easy to see that

$$\frac{1}{\beta_2} \sum_{i=1}^{n} |x_i(t) - y_i(t)| \leq V(x(t)y(t))$$

$$\leq \frac{1}{\beta_1} \sum_{i=1}^{n} |x_i(t) - y_i(t)|, \; t \geq t_0 \in \text{R}. \tag{3.9}$$

Calculating the right derivative $D^+V$ of $V$ and using (1.1) and (3.4) we derive (after simplification) that

$$D^+V\left(x(t), y(t)\right) \leq -\mu \sum_{i=1}^{n} |x_i(t) - y_i(t)| + \sum_{i=1}^{n} |b_i(t) - B_i(t)|$$

$$+\beta_2 \sum_{j=1}^{n} \sum_{i=1}^{n} |a_{ij}(t) - A_{ij}(t)|; \quad t \geq t_0,$$

$$\leq -\mu\, \beta_1\, V\left(x(t), y(t)\right) + \sum_{i=1}^{n} |b_i(t) - B_i(t)|$$

$$+ \beta_2 \sum_{i=1}^{n} \sum_{j=1}^{n} |a_{ij}(t) - A_{ij}(t)|; \quad t \geq t_0, \tag{3.10}$$

An integration of (3.9) over $[t_0, t]$ with an application of (3.8) leads to

$$\frac{1}{\beta_2}\sum_{i=1}^{n}|x_i(t)-y_i(t)| \leq V(x(t),y(t)) \leq V(x(t_0),y(t_0)) + \left[\sum_{i=1}^{n}\sup_{t\in R}|b_i(t)-B_i(t)| + \beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}|a_{ij}(t)-A_{ij}(t)|\right]\frac{1}{\mu\beta_1},$$

$$t \geq t_0,$$

And hence

$$\sum_{i=1}^{n}|x_1(t)-y_1(t)| \leq \left(\frac{\beta_2}{\beta_1}\right)\sum_{i=1}^{n}|x_1(t_0)-y_1(t_0)| + \frac{\beta_2}{\mu\beta_2}\left[\sum_{i=1}^{n}\sup_{t\in R}|b_i(t)-B_i(t)|\right.$$

The construction reason is because is criticized in a number of usage guides: The reason for the long delays was because the costs greatly exceeded the original estimates. One objection to this construction is based on its redundancy: the word because (literally, by cause) contains within it the meaning of reason; thus saying the reason is because is like saying "The cause is by cause," which would never be said. A second objection is based on the claim that because can introduce only adverbial clauses and that reason is requires completion by a noun clause. Critics would substitute that for because in the offending construction: The reason for the long delays in completing the project was that the costs.

$$\left.+\beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}|a_{ij}(t)-A_{ij}(t)|\right]. \qquad (3.11)$$

Now for any $\varepsilon > 0$ if one chooses a $\delta > 0$ such that

$$\delta < \min\left\{\frac{\varepsilon}{3}\frac{\beta_1}{\beta_2}, \frac{\varepsilon}{3}\frac{\beta_1}{\beta^2}\mu\right\} \qquad (3.12)$$

Then (3..6) will be satisfied and this completes the proof

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# . Jia Li's Prescientiial Dynamics And Contential Staticity Of Accentuation Attrition Models Applied To Cosmological And Astrophysical And Quantum Mechanical Engineering Problems

**Abstract**: **Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics ET AL ). We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems. Predator prey model literally form bastion, pillar, post, stylobate, sentinel towards the end of dismasting many of the asymptotic and solutional behaviour in finance, economic, and as we shall see in cosmology, quantum computation and other fields.**

## MODEL FORMULATION

Let $p_1(t, a_1)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_1}{dt} = -b_1 Q_1 + g_1(g_1)Q_1$$

Where $b$ is the natural death rate in the absence of credit types and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_1(0) = 0, \quad g_1(P_1) > 0, \ P_1 \ 0, \qquad \lim_{P_1 \to \infty} g_1(P_1) = (K_1)_1 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_1 + (\rho_a)_1 + D_1(a_1, P_1, Q_1)\rho_1 + 0, \qquad 0 < a_1 \ \infty,$$

$$\rho_1(t, 0) = \int_0^{\infty} n_1(a_1, P_1, Q_1)\rho_1(t_1, a_1)da_1$$

Where $D_1$ is the mortality function and $n_1$ is the fertility function.

Put $_1$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_1}{dt}\right]_1 = x_1\left[\frac{dQ_2}{dt}\right]_1 - x_2\left[\frac{dQ_1}{dt}\right]_1$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_1 - x_4\left[\frac{dQ_2}{dt}\right]_1$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_1 - x_6\left[\frac{dQ_3}{dt}\right]_1$$

Where $x_1, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_1$ and $g_1(P_1)Q_1$

Now when you add the term $(b_1Q_1)$ with $g_1(P_1)(Q_1)$ we are infect see is that $(b_1Q_1)$ is credited incrementally part of $g_1(P_1)(Q_1)$ or instant nearly, when subtracted from $\frac{dQ_1}{dt}$, it means that $g_1(P_1)(Q_1)$ is slowly eaten by $\frac{dQ_1}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_1 = \mu_1(a_1) + (f_1)_1(P_1)_1 + (f_2)_1(Q)_1 \quad (2)$$

$$n_1 = \beta_1(a_1)h_1(Q_1), \quad (3)$$

Where $\mu_1$the age- dependent debit rate whit is $\mu_1(a_1)$ 0 $\forall a > 0$; $(f_1)_1$ is the density- dependent death rate with $f_1(0) = 0$, $(f_1^1)_1(P_1)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_1)_1(P_1) = (k_1)_1$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_1(0) = 0, (f_2^1)_1(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_1(Q_1) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^1(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_1 + (\rho_a)_1 = -\mu_1(a_1) + (f_1)_1P_1) + \big(f_2)_1(Q_1)\big)(\rho_1)_1 \quad (4)$$

$$\rho_1(t_1, 0) = h_1(Q_1)\int_0^\infty \beta_1(a_1)\rho_1(t, a_1)da, \quad (5)$$

$$\frac{dQ_1}{dt}$$

$$= -b_1Q_1 + g_1(P_1)Q_1 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_1$ and $Q_1$ of (4-6) satisfy

$$0 \leq \rho_1(t, a_1) \leq \rho_1^*(t, a_1),$$

$$0 \leq Q_1(t) \leq Q_1^*(t),$$

Where $\rho_1^*(t, a_1)$ and $Q_1^*(t)$, are solutions of

$$(\rho_t^*)_1 + (\rho_a^*)_1 = -(\mu_1(a_1) + (f_1)_1(P_1^*) + (f_2)_1 Q_1^*))\, \rho_1^*, \quad (7)$$

$$\rho_1^*(t, 0) = \int_0^\infty \beta_1(a_1)\, \rho_1^*(t, a_1) da_1 \quad (8)$$

$$\frac{dQ_1^*}{dt} = -b_1 Q_1^* + g_1(P_1^*) Q_1^*, \quad (9)$$

With $P_1^*(t) = \int_0^\infty \rho_1^*(t, a_1) da_1$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_1^*(t)}{dt} = P_1^*(t)\big(r_1 - (f_1)(P_1^*) - (f_2)_1(Q^*)\big), \quad (10)$$

$$\frac{dQ_1^*(t)}{dt} = Q_1^*(t)(-b_1 + g_1(P_1^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_1(a) e^{-\int_0^a \mu_1(\tau) dt} e^{-ra}\, da_1 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_1^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(t)dt}\, da_1 > 1, \qquad (H1)$$

Which guarantees $r_1 > 0$.

*Remark.* If $\beta_1(a) = [\beta_0]_1 a_1 e^{-\alpha_1 a_1}, \mu_1(a_1) \equiv \mu_1$ then the solution of (12) is

$$r_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_1]_1(P_1) = 0, [f_2](Q_1) = \lambda_1 Q_1\ h_1(Q_1) = 1, \qquad and\ \ g_1(P_1) = c_1 P_1$$

$$[\rho_1]_1 + [\rho_a]_1 = -(\mu_1 + \lambda_1 Q_1)\rho_1$$

$$\rho_1(t,0) = \int_0^{\infty} [\beta_0]_1 a_1 e^{-\alpha_1 a_1} \rho_1(t, a_1)\, da_1$$

$$\frac{dQ_1}{dt} = -b_1 Q_1 + c_1 P_1 Q_1$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_1}{dt} = \omega_1 P_1 - \lambda_1 P_1 Q_1$$

$$\frac{dQ_1}{dt} = b_1 Q_1 - c_1 P_1 Q_1$$

Where $\omega_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$ THE EQUILIBRIUM $(\rho_0)_1(a_1), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_1(a_1), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_1 + (\rho_a)_1 = -\big(\mu_1(a_1) + (f_1)_1(P_1)\big)\rho,$$

$$\rho_1(t,0) = \int_0^\infty \beta_1(a_1)\rho_1(t,a_1)da_1 \qquad (13)$$

The solution of (13) has the form of

$$\rho_1(t,a_1) = Ce_1^{r(t-a_1)} - \int_0^a (\tau)d\tau - \int_0^1 (f_1)_1\big(P_1(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_1(t) = C_1 e^{r_1 t - \int_0^t (f_1)_1(P_1(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_1^{-1}(r)$ if $k_1 > r$.

This is equivalent to

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{(k_1)_1 a_1} da_1 < 1.$$

In fact, if we define

$$R_1(s_1) = \int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau)d\tau} e^{-s_1 a_1}\, da_1$$

$R(s)$ is a decreasing function. $R_1(r_1) = 1$ And $R_1(k_1)_1 < 1$ imply $r_1 < (k_1)_1$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_1(a_1), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{k_1 a_1} da_1 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_1 \equiv \int_0^\infty (\rho_0)_{11}(a_1)da_1$ *satisfies*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{-(f_1)_1(P_0)a_1} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_1(a_1), 0)$

Let $u_1(t, a_1) = \rho(t, a_1) - \rho_0(a_1)$ Then

$$N_1(t) \equiv \int_0^\infty u_1(t, a_1)da_1 = \int_0^\infty \rho_1(t, a_1)da_1 - \int_0^\infty (\rho_0)_1(a_1)da = P_1(t) - P_0)_1$$

And

$$(u_1)_1 + (u_a)_1 = (\rho_t)_1 + (\rho_a)_1 - (\rho_0^1)_1(a_1)$$

$$= -\mu_1(a_1) + (f_1)_1(P_1) + (f_2)_1(Q_1))(u_1 + (\rho_0)_1) + \mu_1(a_1)(\rho_0)_1 + (f_1)_1(P_0)_1(\rho_0)_1$$

$$= -\mu_1(a_1)u_1 - (f_1)_1(P_1) + (f_2)_1(Q_1))u_1 - (\rho_0)_1(f_1)_1(P_1) - (f_1)_1(P_0)_1$$

$$+ (f_2)_1(Q_1)), \qquad (15)$$

Where $(\rho_0^1)_1(a_1)$ denotes $d(\rho_0)_1(a_1)/da_1$

Expansion $(f_1)_1(P_1)$ and $(f_2)_1(Q_1)$ at $(P_0)_1$ and 0 yield respectively

$$(f_1)_1(P_1) = (f_1)_1((P_0)_1 + N_1)$$

$$= (f_1)_1(P_0)_1 + (f_1^1)_1(P_0)_1N_1 + o_1(N_1), \qquad (16)$$

$$(f_2)_1(Q_1)$$

$$= (f_2^1)_1(0)Q + o_1(Q_1) \qquad (17)$$

Write $f_1(P_0) = f_{10}, f_1'(P_0) = f_{10}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_1)_1 + (u_a)_1 \approx -(\mu_1(a_1) + (f_{10})_1)u - (f_{10}^1)_1(\rho_0)_1(N_1) - (f_{20}^1)_1(\rho_0)_1(Q_1)$$

$$u_1(t,0) = \rho_1(t,0)(\rho_0)_1(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t,a) + (\rho_0)_1(a_1)da_1 - (\rho_0)_1(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t,a_1)da_1 + (h_1(Q_1) - 1)(\rho_0)_1(0) \qquad (18)$$

Expanding

$$h_1(Q_1) = 1 + (h^1)_1(0)Q_1 + 0(Q_1) \equiv 1 + h_0'Q_1 + 0Q_1$$

And substituting it into (18), we arrive at

$$u_1(t,0) \approx \int_0^\infty \beta_1(a_1)u_1(t,a_1)da_1 + (h_{a_1}')(\rho_0)_1Q_1$$

The linearization is the following:

$$[u_1] + [u_1] = -(\mu_1(a_1) + [t_{10}]_1u_1 - [\alpha_1]_1(a_1)N_1 - [\alpha_2]_1(a_1)Q_1$$

$$u_1(t,0) = \int_0^\infty \beta_1(a_1)u_1(t,a_1)da_1 + [\xi_1]Q_1 \qquad (19)$$

$$\frac{dQ_1}{dt} = -b_1Q_1$$

where

$$[\alpha_1]_1(a_1) = [t_1^1 a]_1[p, o]_1(a_1)d2(a_1) = [f_{20}^1]_1(\rho\, o\,)_1(a_1)\xi = [\mathrm{h\ o}^1]_1[\rho \mathrm{o}]_1(\mathrm{o}) \text{ and } \mathrm{b}_1 = \mathrm{b}_1 - g_1(P_1)$$

If $g_1(P_0) > b$, clearly, $(\rho_0)_1(a_1), 0)$ is unstable and the predator population is persistent. Assume $g_1(P_0)_1 < b$.

Then, $Q_1(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_1 + u_a = -(\mu_1(a_1) + (f_{10})_1)u_1 - (\alpha_1)_1(a_1)N_1$$

$$= -(\mu_1(a_1) + f_{10})u_1 - \alpha_1(a_1)\int_0^\infty u_1(t, a_1)da_1$$

$$u_1(t, 0) = \int_0^\infty \beta_1\,(a_1)u_1(t, a_1)da_1$$

Using the method developed by Cushing [2, 3], we suppose

$$u_1(t, a_1) = \tilde{\mathrm{u}}_1(a_1)^{e^{\gamma(t-a_1)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_1}\, da_1$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_1)$ gives,

$$\tilde{\mathrm{u}}_1(a_1) = e^{-\int_0^a \mu_1(\tau_1)d\tau - (f_{10})_1 a}\left(\tilde{\mathrm{u}}_1(0) - w\int_0^a e^{\int_0^\tau \mu_1(s_1)ds + (f_{10})}\tau_1\,(\alpha_1)_1(\tau_1)e^{\gamma_1\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_1(a_1)e^{-\gamma_1 a_1 - (f_{10})_1 a_1 - \int_0^a \mu_1(\tau_1)d\tau}\right)\tilde{\mathrm{u}}_1(0)$$

$$+w\int_0^{\infty}\beta_1(a_1)e^{-\gamma_1 a_1-(f_{10})_1 a_1-\int_0^a \mu_1(\tau_1)d\tau}\int_0^a(\alpha_1)_1(\tau_1)e^{-\gamma_1\tau_1}e^{(f_{10})_1\tau_1+\int_0^{\tau}\mu_1(s_1)ds}d\tau_1\,da_1$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty}e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}da\right)u(0)$$

$$+\left(-1-\int_0^{\infty}e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}\int_0^a e^{f_{10}\tau+\int_0^{\tau}\mu(s)ds+y\tau}\alpha_1(\tau)d\tau\,da\right)w=0.\ (23)$$

Define $Y(\gamma,a)=e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta\,(a)Y(\gamma,a)da\right)u(0)$$

$$+w\int_0^{\infty}\beta(a)Y\,(\gamma,a)\int_0^a\alpha_1(\tau)Y^1(\gamma,\tau)d\tau\,da=0,(24)$$

$$\left(\int_0^{\infty}Y(\gamma,a)da\right)u(0)$$

$$+\left(-1-\int_0^{\infty}Y(\gamma,a)\int_0^a\alpha_1(\tau)Y^{-1}(\gamma,\tau)d\tau\,da\right)w=0.$$

Equation (20) has a solution of the form of $u(t,a)=u(a)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $u(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I(y,a)=\int_0^a\rho_0(\tau)Y^{-1}\,(\gamma,\tau)d\tau$. We obtain the characteristic equation

$$f_{10}^{l}\left(\int_0^{\infty} Y(\gamma,a)da\right)\left(\int_0^{\infty}\beta(a)Y(\gamma,a)I(\gamma,a)da\right)$$

$$+\left(1-\int_0^{\infty}\beta(a)Y(\gamma,a)da\right)\left(1+f_{10}^{l}\int_0^{\infty}Y(\gamma,a)I(\gamma,a)da\right)=0. \quad (25)$$

Since

$$\rho_0(a)=\rho_0(0)e^{-\int_0^a \mu(\tau)d\tau - f_{10}a}$$

$$I(\gamma,a)=\int_0^a \rho_0(\tau)e^{\gamma\tau+f_{10}\tau+\int_0^{\tau}\mu(s)ds}d\tau$$

$$=\frac{\rho_0(0)}{\gamma}(e^{\gamma a}-1)$$

Let

$$E(a)=e^{-f_{10}a-\int_0^a \mu(\tau)d\tau}$$

Then $[Y_1]_1[\gamma_1]_1,a_1)=\ E_1(a_1)e^{-\gamma_1 a_1}$. by a simple calculation,

$$\int_0^{\infty}\beta_1(a_1)Y_1(\gamma_1,a_1)I(\gamma_1,a_1)da=\frac{[\rho_0]_1(0)}{\gamma}\left(1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}\ da\right);$$

$$\int_0^{\infty}Y_1(\gamma_1,a_1)I(\gamma_1,a_1)da=\frac{[\rho_0]_1(0)}{\gamma_1}\left(\int_0^{\infty}E_1(a_1)e^{-\gamma_1 a_1}\ da_1\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_1(0)}{\gamma_1}[f_{10}^{1}]\int_0^{\infty}E_1(a_1)e^{-\gamma_1 a_1}da_1\ (\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1)$$

$$+\left(1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right).$$

$$\left(1+\frac{(\rho_0)_1(0)}{\gamma_1}(f_{10}^{'})_1\left(\int_0^{\infty}E_1(a_1)da_1-\int_0^{\infty}E_1(a_1)e^{-\gamma_1 a_1}da_1\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right)(1+\frac{(f_{10}^{'})_1}{\gamma_1}(\rho_0)_1(0)\int_0^{\infty}E_1(a_1)da_1)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_1(a_1)Y_1(\gamma_1 a_1)da_1=1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\neq 0,$$

Since $\int_0^{\infty}\beta_1(a_1)Y_1(\gamma_1 a_1)\int_0^{a}\alpha_1(\tau_1)Y^{-1}(\gamma_1,\tau_1)d\tau_1 da_1>0$. Therefore, the characteristic solution is

$$\gamma_1=-(f_{10}^{'})_1(\rho_0)_1\int_0^{\infty}E_1(a_1)da_1.$$

Since $f_1'(P_0)_1>0$, we obtain the following result

THEOREM 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_1),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_1),0)$ which is related to the extinction of debit summation as has been referred to in the foregoing and to The possible persistence of the prey

population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_1 + (\mu_1(a_1) + [f_1]_1(P_1^*) + [f_2]_1(Q_1^*))P_1^* = 0, \quad (26)$$

$$p_1^*(0) = h_1(Q_1^*)\int_0^\infty \beta_1(a_1)p_1^*(a_J)p_1^*(a_1)da_1 \quad (27)$$

$$g_1(P_1^*) - b_1 = 0, \quad (28)$$

Where $P_1^* = \int_0^\infty p_1^*(a_1)\, da_1$

Equation (28) yields $P^* = g^{-1}(b)$,

From (26),

$$p_1^*(a) = p_1^*(0)e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}$$

Substituting into (27), we need to solve

$$\left(1 - h_1(Q_1^*)\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}da_1\right)p_1^*(0) = 0$$

For $Q_1^*$

Fix $p_1^*(0)$ and $P_1^*$ Let

$$H_1(z_1) = 1 - h_1(z_1)\int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1}da_{,1}$$

And assume that

$$[k_3]_1 \int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(g_1^{-1}(b_1))a_1 - [k_2]_1 a_1}\, da_1 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty} [H_1(z_1)] > 0.$$

Clearly,

$$\frac{dH_1}{dz_1} = h_1(z)(\,[f_2^{'}]_1(z_1)\,) \int_0^{\infty} a_1\beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - f_1(P_1^*)a_1 - f_2(z_1)a_1} da_1$$

$-h^{'}(z_1) \int_0^{\infty} \beta_1(a_1)\, e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1}\ da_1$

By the assumptions on $\left[f_2^1\right]_1 and\ [h^1]^1 dH_1/dz > 0$. Hence, if

$$H_1(0) = 1 - \int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1} da_1 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_1(0) < 0$ follows if $g_1^-(b_1) < [f_1^{-1}]_1(r_1)$, where $r_1$ is the solution of (12).

Indeed, since

$$R_1(s_1) = \int_0^{\infty} \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 e^{-y\frac{a}{1}}} da_1$$

Is a decreasing function of $s_1$ when $P_1^* = [g^{-1}]_1(b_1) < [f_1^{-1}]_1(r_1)$, that is $[f_1]_1(P_1^*) < r$, $R\big(f_1(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_1(g^{-1}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1\left(g_1^{-1}(b)\right)a_1} > 1. \qquad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_1^*(a_1), Q_1^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_1^*(a_1), Q_1^*)$ is more complicated. We again consider perturbations $m_1(t_1 a_1) = p_1(t_1, a_1) - p_1^*(a_1)$ and $W_1(t) = Q_1(t) - Q^*$.Then

$M_1(t) \equiv \int_0^\infty m_1(t, a_1)da_1 = P_1(t) - P_1^*,$

$$[m_1]_1 + [m_a]_1 = [p_1]_1 + [p_a]_1 - \left[p^*\right]_1(a_1) \approx -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1$$

$$-[f_{1*}^1]_1 p_1^*(a_1)M - [f_{2*}^1]_1 p_1^*(\mathrm{a})W_1,$$

$$m_1(t,0)p_1(t,0) - p_1^*(0) \approx h_* \int_0^\infty \beta_1(a_1)m_1(t,a_1)\,da + \frac{[h_*^1]_1}{[h_*]_1}\, p_1^*(0)W_1$$

$$\frac{dW_1}{dt} = \frac{dQ_1}{dt} \approx [g_*^1]_1 Q_1^* M_{,1}$$

Where $[f_{1*}]_1$ denoting $[f_1(P^*)]_1$, $[f_{1*}]_1$ denoting $[f_2(P^*)]_1$, $\left[f_{1*}^{'}\right]_1$ denoting $\left[f^{'}{}_1(P^*)\right]_1 \left[f_{2*}^{'}\right]_1$ denoting

$[f_2]_1\,(Q_1^*,\quad h_{1*}^{'}$ denoting $h_*^{'}$ denoting $h_1^{'}\,(Q_1^*)$, and $\left[g_*^{'}\right]_1$

Denoting $\left[g^{'}\right]_1(P_1^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_1]_1 + [m_a]_1 = -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1 - [f'_{1*}]_1[p_1^*](a_1)M_1 - [f'_{2*}]_1 p_1^*(a_1)W_1$$

$$m_1(t,0) = [h_*]_1 \int_0^\infty \beta_1(a_1)m_1(t,a_1)da_1 + \frac{[h_*^1]_1}{[h^*]_1}[p^*]_1(0)\,W_{,1} \qquad (30)$$

$$\frac{dW_1}{dt} = [g_*^1]_1[Q_1^*]M_1$$

Integrating $[m_1]_1(t,a_1)$ with respect to a from 0 to $\infty$, we have

$$\frac{dM_1(t)}{dt} = \int_0^\infty [mt]_1(t,a_1)da_1$$

$$= -\int_0^\infty [m_a]_1\,(t,a_1)da_1 - \int_0^\infty \mu(a_1)m(t,a_1)da_1$$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f'_{1*}]_1 \int_0^\infty p_1^*(a_1)da_1\right)M_1 - [f_{2*}^1]_1 \int_0^\infty p_1^*(a_1)da_1 W_1$$

$$= m_1(t,0) - \int_0^\infty \mu_1(a_1)\,m_1(t,a_1)da_1$$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f_{1*}']_1 P_1^*\right)M_1 - [f_{2*}']_1 P_1^* W_{,1}$$

From (30), we arrive at the equations

$$\frac{dW_1(t)}{dt} = \int_0^\infty ([h_*]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1$$

$$-\left([f_{1*}] + [f_{2*}]_1 + [f_{1*}']_1 P_1^*\right)M_1 + \left(\frac{[h_*']_1}{[h_*]^1}\,P_1^*(0) - f_{2*}' P_1^*\right)W_1,$$

(31)

$$\frac{dW_1(t)}{dt}\left[g^{'}{}_{*}\right]_1 Q_1^{*} M_{.1}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{*}]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1 = ([h_{*}]_1\beta_1(\theta_1) - \mu_1(\theta_1))M_{,1}$$

Where $0< \theta_1 = \theta_1(t) < \infty$.

And if $a_1^2\,[a_1^2]_1<4\,[a_2]_1<4\,a_2 e$,

Re $\lambda_1\left[\lambda_1\right]_1 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_1 = -\frac{1}{2}a_1[a2\ ]_1 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_{*}\,[h_{*}]_1\beta_1(a_1) - \mu_1(a_1)\right|$. Then$|n_1(t)|$ ⊠Δ∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium($\rho^{*}(a), \boldsymbol{Q}^{*}$) is asymptotically stable if either $a_1^2 <4\,a_2 e$ and $[a_1] > 2\Delta$,or $[a_1^2]$ ⊠4$[a_2]$e and $(a_1 - \sqrt{(a_1^2)_1 - 4(a_2)_1 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+1}\right)$ and $\beta_1(a_1) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_1(a_1)\mathrm{e}^{-\int_0^{\mathrm{a}} \mu_1(\tau_1)+d\tau_1 \mathrm{da}} = \frac{1}{(\mathrm{m}+\mu_0+1)^2} + \frac{2}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_1 + [\mu_0]$ 0.5. Assume$[f_1](x_1)= K_1 x_1$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_1(x_1)= [0x_1]/(1 + x_1)$; the functional responses $[[f_2](x) = [0]\, x_1/(1 + x_1)$and $h_1(x_1)=(1+\omega_1 x_1)/(1 + x_1)$(See Holling [7] and Rashevsky [13])

Equation (14),now has the form of

$$\frac{m_1 + [\mu_0] + [\mathrm{K}][P_0])}{[m_1]_1 + (\mu_0)_1 + [\mathrm{K}][P_0])_1^3\ [\mathrm{K}]_1[P_0]_1^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_1$i e $[P_0] < b/(\delta_1 - b_1)$then the equilibrium$(P_0)\ [P_0]_1(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_1$

$[f_{1*}] = K_1\ \mathrm{P}^*\ [f_{2*}] = \frac{\sigma \mathrm{Q}_1^*}{1+\mathrm{Q}_1^*} \quad [h_*] = \frac{1+\omega_1 \mathrm{Q}_1^*}{1+\mathrm{Q}_1^*}$

$[f_{1*}] = K_1\ [\mathrm{f}_{2*}^1] = \frac{\sigma_1}{(1+\mathrm{Q}_1^*)^2}[\ \mathrm{g}_*^1] = \frac{\sigma_1}{(1+\mathrm{P}_1^*)^2} \quad [\mathrm{h}_*^1] = \frac{\omega_1 - 1}{(1+\mathrm{Q}_1^*)^2}$

## AN ANALOGY BETWEEN PREDATOR PREY AND ACCENTUATION ATTRION MODELS

**Best example how the model could be applied is the Schrödinger's equation itself wherein RHS gobbles up LHS, and the terms on RHS do the same with each other. It is like adding milk to water. Whether it happens instantaneously or time lag is there is a different question altogether dealt at the opportune points**

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| **Best example that can be given is the formulation of the quantum wave equation albeit in a nascent form for the recently found gravitational waves. This has been studied in length in various papers of gravitational waves, wherein the attribution of how the model is used is described. Note the calculation holds for only one module. This is in consideration to the fact that the relation is beteen the both debit systems and credit systems which as we have enucleated in the foregoing bears ample testimony for the application of the model** |
| **Gravitational quantum wave equation:** |
| **Equation Time-dependent equation (Wikipedia)** |
| **(1) In quantum mechanics, the Schrödinger equation is (=) a partial differential equation that describes how the quantum state of a quantum system changes with (e&eb) time. It was formulated in late 1925, and published in 1926, by the Austrian physicist Erwin Schrödinger.[1]** |
| **(2) In classical mechanics Newton's second law, (F = ma), is used to make (eb) a mathematical prediction as to what path a given system will take following a set of known initial conditions.** |
| **(3) In quantum mechanics, the analogue of Newton's law is (=) Schrödinger's equation for (e) a quantum system (usually atoms, molecules, and subatomic particles whether free, bound, or localised).** |
| **(4) It is not a simple algebraic equation, but in general a linear partial differential equation, describing (eb) the time-evolution of the system's wave function (also called a "state function").[2]:1–2** |
| **(5) The concept of a wavefunction is (=) a fundamental postulate of quantum mechanics.** |
| **(6) Although Schrödinger's equation is often presented as a separate postulate, some authors [3]: Chapter 3 show that some properties resulting from Schrödinger's equation may be deduced just from (e) symmetry principles alone, for example the commutation relations.** |
| **(7) Generally, "derivations" of the Schrödinger equation demonstrate (eb) its mathematical plausibility for describing (eb) wave–particle duality, but to date there are no universally accepted derivations of (e) Schrödinger's equation from (e) appropriate axioms.** |
| **(8) In the Copenhagen interpretation of quantum mechanics, the wave function is (=) the most complete description that can be given of (e) a physical system.** |
| **(9) Solutions to Schrödinger's equation describe (eb) not only molecular, atomic, and subatomic systems, but also macroscopic systems, possibly even the whole universe.[4]:292ff** |

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| **(10) The Schrödinger equation is consistent with (e&eb) classical mechanics but not with (e&eb) special relativity.** |
| **(11) Making quantum mechanics consistent with (e&eb) special relativity is one of the great achievements of (e) Quantum Field Theory and requires (e) the introduction of the creation and annihilation operators so that (e) the number of particles can no longer be considered constant.** |
| **(12) The Schrödinger equation is not the only way to make predictions in (eb) quantum mechanics—other formulations can be used, such as Werner Heisenberg's matrix mechanics, and Richard Feynman's path integral formulation** |
| **(13) The form of the Schrödinger equation depends on (=) the physical situation (see below for special cases).** |
| **(14) The most general form is the time-dependent Schrödinger equation, which gives a description of a system evolving with time: [5]:143** |
| **(15) A wave function that satisfies the non-relativistic Schrödinger equation with V = 0. In other words, this corresponds to (e &eb) a particle traveling freely through (e&eb) empty space. Time-dependent** |
| **Schrödinger equation (general)** |
| **(16) i ℏ ∂ ∂ t Ψ ( r , t ) = H ^ Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t) where i is the imaginary unit, ħ is the Planck constant divided by 2π, the symbol ∂/∂t indicates a partial derivative with respect to time t, Ψ (the Greek letter psi) is the wave function of (e) the quantum system, and Ĥ is the Hamiltonian operator (which characterizes the total energy of any given wave function and takes different forms depending on the situation).** |
| **(17) Each of these three rows is a wave function which satisfies (eb) the time-dependent Schrödinger equation for a harmonic oscillator. Left: The real part (blue) and imaginary part (red) of the wave function. Right:** |
| **(18) The probability distribution of finding the particle with (e&eb) this wave function at a given position.** |
| **(19) The top two rows are examples of stationary states, which correspond to (e&eb) standing waves.** |
| **(20) The bottom row is an example of a state which is (=) not a stationary state.** |
| **(21) The right column illustrates (eb) why stationary states are (=) called "stationary".** |

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| **(22) The most famous example is the non-relativistic Schrödinger equation for (e) a single particle moving in an electric field (but not a magnetic field; see the Pauli equation): [6]** |
| **Time-dependent Schrödinger equation** |
| **(Single non-relativistic particle)** |
| **(23) i ℏ ∂ ∂ t Ψ ( r , t ) = [ − ℏ 2 2 μ ∇ 2 + V ( r , t ) ] Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)** |
| **(24) Here μ is (=) the particle's "reduced mass", V is its potential energy, ∇2 is the Laplacian (a differential operator), and Ψ is (=) the wave function (more precisely, in this context, it is called the (=) "position-space wave function").** |
| **(25) In plain language, it means "total energy equals (=) kinetic energy plus potential energy", but the terms take unfamiliar forms for (e) reasons explained below.** |
| **(26) Given the particular differential operators involved, this is (=) a linear partial differential equation.** |
| **(27) It is also a diffusion equation, but unlike the heat equation, this one is (=) also a wave equation given (e&eb) the imaginary unit present in the transient term.** |
| **(28) The term "Schrödinger equation" can refer to (e) both the general equation (first box above), or the specific nonrelativistic version (second box above and variations thereof).** |
| **(29) The general equation is indeed quite general, used throughout quantum mechanics, for everything from the Dirac equation to quantum field theory, by plugging in various complicated expressions for the Hamiltonian. The specific nonrelativistic version is a simplified approximation to reality, which is quite accurate in many situations, but very inaccurate in others (see relativistic quantum mechanics and relativistic quantum field theory)** |
| **(30) To apply the Schrödinger equation, the Hamiltonian operator is set up for the system, accounting for the kinetic and potential energy of the particles constituting the system, then inserted into the Schrödinger equation. The resulting partial differential equation is solved for the wave function, which contains information about the system.** |
| **Time-independent equation** |

**(31) The time-dependent Schrödinger equation described above predicts that wave functions can form standing waves, called stationary states (also called "orbitals", as in atomic orbitals or molecular orbitals). These states are important in their own right, and if the stationary states are classified and understood, then it becomes easier to solve the time-dependent Schrödinger equation for any state. Stationary states can also be described by a simpler form of the Schrödinger equation, the time-independent Schrödinger equation. (This is only used when the Hamiltonian itself is not dependent on time explicitly. However, even in this case the total wave function still has a time dependency.)**

**Time-independent Schrödinger equation (general)**

**(32) E Ψ = H ^ ⊠ Ψ {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi } {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi }**

**In words, the equation states:**

**(33) When the Hamiltonian operator acts on a certain wave function Ψ, and the result is proportional to the same wave function Ψ, then Ψ is a stationary state, and the proportionality constant, E, is the energy of the state Ψ.**

**(34) The time-independent Schrödinger equation is discussed further below. In linear algebra terminology, this equation is an eigenvalue equation.**

**(35) As before, the most famous manifestation is the non-relativistic Schrödinger equation for a single particle moving in an electric field (but not a magnetic field):**

**Time-independent Schrödinger equation (single non-relativistic particle)**

**(36) E Ψ ( r ) = [ − ℏ 2 2 μ ∇ 2 + V ( r ) ] Ψ ( r ) {\displaystyle E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )} E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )**

**Gravitational waves: (Wikipedia)**

**(37) In theory, the loss of energy through gravitational radiation could eventually drop the Earth into the Sun. However, the total energy of the Earth orbiting the Sun (kinetic energy + gravitational potential energy) is about 1.14×1036 joules of which only 200 Watts (joules per second) is lost through gravitational radiation, leading to a decay in the orbit by about 1×10−15 meters per day or roughly the diameter of a proton. At this rate, it would take the Earth approximately 1×1013 times more than the current age of the Universe to spiral onto the Sun. This estimate overlooks the decrease in r over time, but the majority of the time the bodies are far apart and only radiating slowly, so the difference is unimportant in this example.**

**More generally, the rate of orbital decay can be approximated by[39]**

**(38) d r d t = − 64 5 G 3 c 5 ( m 1 m 2 ) ( m 1 + m 2 ) r 3 , {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,} {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,}**

**where r is the separation between the bodies, t time, G Newton's constant, c the speed of light, and m1 and m2 the masses of the bodies. This leads to an expected time to merger of [39]**

**t = 5 256 c 5 G 3 r 4 ( m 1 m 2 ) ( m 1 + m 2 ) . {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.} {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.}**

**(39) For example, a pair of solar mass neutron stars in a circular orbit at a separation of 1.89×108 m (189,000 km) has an orbital period of 1,000 seconds, and an expected lifetime of 1.30×1013 seconds or about 414,000 years. Such a system could be observed by LISA if it were not too far away. A far greater number of white dwarf binaries exist with orbital periods in this range. White dwarf binaries have masses in the order of the Sun, and diameters in the order of the Earth. They cannot get much closer together than 10,000 km before they will merge and explode in a supernova which would also end the emission of gravitational waves. Until then, their gravitational radiation would be comparable to that of a neutron star binary.**

**(40) By substituting value of dr/dt in Schrödinger's wave equation we get the quantum wave equation for gravitational waves.**

In order to have (H3) and (H4) hold, $m, \mu_0$, K, b, δ, and $\omega$ need to satisfy

$$\frac{\omega_1}{(m+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1+b_1}+\sigma)^2}+\frac{2\omega_1}{([m_1]_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1+b_1}+\sigma_1)^3}<1$$

And

$$\frac{1}{(m_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1-b_1})^2}+\frac{2\omega_1}{(m_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1-b_1})^3}>1$$

When the inequalities (33) and (34) hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$(m_1 + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 - b_1} + (m_1 + [\mu_0]_1 + \frac{K_1 b_1}{\delta_1 - b_1} + \sigma_1)\, Q_1^*)^3$

$= m_1 + [\mu_0]_1 + 3 + \frac{K_1 b_1}{\delta_1 - b_1}(1 + Q^*)(1 + \omega_1 Q_1^*$

Now

$$\left[a_1^2\right] - 4\left[a_2\right] e = \left(\frac{2K_1 b_1}{\delta_1 - b_1}\right)^2 + \frac{4K_1 b_1 \sigma_1 Q_1^*}{(\delta_1 - b_1)(1 + Q_1^*)} + \frac{\sigma_1^2 Q_1^{*2}}{(1 + Q_1^*)^2}$$

$$\frac{4\, b_1 \sigma_1 (\delta_1 - b) Q^*}{\delta_1 (1 + Q_1^*)^2} - \frac{4\, (\delta_1 - b_1)^2 (1 - \omega_1) Q_1^*}{\delta_1 (1 + \omega_1 Q_1^*)(1 + Q_1^*)} P^*(0)$$

And

$\Delta = \lim_{0 < a < \infty} |[h_*]_1\, a e^{mu}_1\ \ [\mu_0]_1 - \frac{a_1}{a_1 + 1}$

Hence if $[a_1^2 > 4(a_2)_1 e$ and

$$[a_1] = \frac{2K_1 b_1}{\delta_1 - b_1} + \frac{\sigma_1 Q_1^*}{1 + Q_1^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_1]_1 = [f_{1*}]_1 + [f_{2*}]_1 + [f_{1*}] P_1^*$$

$$[a_2]_1 = \left[f^{1}_{2*}\right] P^{*}_{1} - \frac{h^{1}_{*}}{h_{*}} P^{*}(0)$$

$$e=[g^{1}_{*}]_1 Q^{*}_{1}$$

$$\eta_1(t) = \left[h_{*}\right]_1 \beta_1 \left(\theta_1(t)\right) - \mu_1(\theta(t)),$$

The system(31)can be written as the following system:

$$[M(t)]\left[a_1\right]M_1 - [a_2]_1 W_1 - [n(t)]M_1$$

$$w_1(t) = eM_1$$

Write

$$A=\begin{pmatrix} -[a_1]_1 - [a_2]_1 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_1 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_1 \\ W_1 \end{pmatrix}$$

The solution of (32)is

$$X_1(t) = e^{A,t}[X]_1(0) + \int_0 e^{At}$$

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# Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

## (Module Three)

**Abstract: Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-435) 1990. We demonstrate its applicability cosmology/ quantum computation economics et al. As stated in the foregoing we shall apply model to various cosmological problem and show the relevance to the same. . To exemplify model can be applied to the following situational significations: Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an**

**outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 Masaru Shibata and Takashi Nakamura** perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. **Authors** use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society **Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995**

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(g_3)Q_3$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_3(0) = 0, \quad g_3(P_3) > 0, \quad P_3 \text{ 0}, \qquad \lim_{P_3 \to \infty} g_3(P_3) = (K_3)_3 \text{ ∞}. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_3 + (\rho_a)_3 + D_3(a_3, P_3, Q_3)\rho_3 + 0, \qquad 0 < a_3 \text{ ∞},$$

$$\rho_3(t,0) = \int_0^\infty n_3(a_3, P_3, Q_3)\rho_3(t_3, a_3)da_3$$

Where $D_3$ is the mortality function and $n_3$ is the fertility function.

Put $_3$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_3}{dt}\right]_3 = x_3\left[\frac{dQ_2}{dt}\right]_3 - x_2\left[\frac{dQ_3}{dt}\right]_3$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_3 - x_4\left[\frac{dQ_2}{dt}\right]_3$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_3 - x_6\left[\frac{dQ_3}{dt}\right]_3$$

Where $x_3, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_3$ and $g_3(P_3)Q_3$

Now when you add the term $(b_3Q_3)$ with $g_3(P_3)(Q_3)$ we are infect see is that $(b_3Q_3)$ is credited incrementally part of $g_3(P_3)(Q_3)$ or instant nearly, when subtracted from $\frac{dQ_3}{dt}$ , it means that $g_3(P_3)(Q_3)$ is slowly eaten by $\frac{dQ_3}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_3 = \mu_3(a_3) + (f_3)_3(P_3)_3 + (f_2)_3(Q)_3 \quad (2)$$

$$n_3 = \beta_3(a_3)h_3(Q_3), \quad (3)$$

Where $\mu_3$the age- dependent debit rate whit is $\mu_3(a_3)$ 0 $\forall a > 0$; $(f_3)_3$ is the density- dependent death rate with $f_3(0) = 0$, $(f_3^3)_3(P_3)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_3)_3(P_3) = (k_3)_3$ ; $f_2$ is the functional response of

creditors to prey with $[f_2]_3(0) = 0, (f_2^3)_3(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_3\ (Q_3) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 3, h^3(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_3 + (\rho_a)_3 = -\mu_3(a_3) + (f_3)_3 P_3) + \big(f_2)_3(Q_3)\big)(\rho_3)_3 \tag{4}$$

$$\rho_3(t_3, 0) = h_3(Q_3)\int_0^\infty \beta_3\,(a_3)\rho_3(t, a_3)da, \tag{5}$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(P_3)Q_3 \tag{6}$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_3$ and $Q_3$ of (4-6) satisfy

$$0 \ \rho_3(t, a_3) \ \rho_3^*(t, a_3),$$

$$0 \ Q_3(t) \ Q_3^*(t),$$

Where $\rho_3^*(t, a_3)$ and $Q_3^*(t)$, are solutions of

$$(\rho_t^*)_3 + (\rho_a^*)_3 = -(\mu_3(a_3) + (f_3)_3(P_3^*) + (f_2)_3 Q_3^*))\,\rho_3^*, \tag{7}$$

$$\rho_3^*\,(t, 0) = \int_0^\infty \beta_3(a_3)\,\rho_3^*\,(t, a_3)da_3 \tag{8}$$

$$\frac{dQ_3^*}{dt} = -\,b_3\,Q_3^* + g_3(P_3^*)Q_3^*, \tag{9}$$

With $P_3^*(t) = \int_0^\infty \rho_3^*\ (t, a_3)da_3$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_3^*(t)}{dt} = P_3^*(t)\big(r_3 - (f_3)(P_3^*) - (f_2)_3(Q^*)\big), \quad (10)$$

$$\frac{dQ_3^*(t)}{dt} = Q_3^*(t)(-b_3 + g_3(P_3^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_3(a)e^{-\int_0^a \mu_3(\tau)dt}e^{-ra}\ da_3 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_3^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(t)dt}\ da_3 > 1, \quad (H1)$$

Which guarantees $r_3 > 0$.

*Remark.* If $\beta_3(a) = [\beta_0]_3 a_3 e^{-\alpha_3 a_3}, \mu_3(a_3) \equiv \mu_3$ then the solution of (12) is

$$r_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_3]_3(P_3) = 0, [f_2](Q_3) = \lambda_3 Q_3\ h_3(Q_3) = 1, \quad and\ \ g_3(P_3) = c_3 P_3$$

$$[\rho_3]_3 + [\rho_a]_3 = -(\mu_3 + \lambda_3 Q_3)\rho_3$$

$$\rho_3(t,0) = \int_0^{\infty} [\beta_0]_3 a_3 e^{-\alpha_3 a_3} \rho_3(t,a_3)\, da_3$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + c_3 P_3 Q_3$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_3}{dt} = \omega_3 P_3 - \lambda_3 P_3 Q_3$$

$$\frac{dQ_3}{dt} = b_3 Q_3 - c_3 P_3 Q_3$$

Where $\omega_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$

## THE EQUILIBRIUM $(\rho_0)_3(a_3), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_3(a_3), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_3 + (\rho_a)_3 = -\big(\mu_3(a_3) + (f_3)_3(P_3)\big)\rho,$$

$$\rho_3(t,0)$$

$$= \int_0^\infty \beta_3(a_3)\rho_3(t,a_3)da_3 \qquad (13)$$

The solution of (13) has the form of

$$\rho_3(t,a_3) = Ce_3^{r(t-a_3)} - \int_0^a (\tau)d\tau - \int_0^3 (f_3)_3\big(P_3(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_3(t) = C_3 e^{r_3 t - \int_0^t (f_3)_3(P_3(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_3^{-3}(r)$ if $k_3 > r$. This is equivalent to

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{(k_3)_3 a_3} da_3 < 1.$$

In fact, if we define

$$R_3(s_3) = \int_0^\infty \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau)d\tau} e^{-s_3 a_3}\, da_3$$

$R(s)$ is a decreasing function. $R_3(r_3) = 1$ And $R_3(k_3)_3 < 1$ imply $r_3 < (k_3)_3$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_3(a_3), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{k_3 a_3} da_3$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_3 \equiv \int_0^\infty (\rho_0)_{33}(a_3)da_3$ *satisfies*

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{-(f_3)_3(P_0)a_3} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_3(a_3), 0)$

Let $u_3(t, a_3) = \rho(t, a_3) - \rho_0(a_3)$ Then

$$N_3(t) \equiv \int_0^{\infty} u_3(t, a_3) da_3 = \int_0^{\infty} \rho_3(t, a_3) da_3 - \int_0^{\infty} (\rho_0)_3(a_3) da = P_3(t) - P_0)_3$$

And

$$(u_3)_3 + (u_a)_3 = (\rho_t)_3 + (\rho_a)_3 - (\rho_0^3)_3(a_3)$$

$$= -\mu_3(a_3) + (f_3)_3(P_3) + (f_2)_3(Q_3))(u_3 + (\rho_0)_3) + \mu_3(a_3)(\rho_0)_3 + (f_3)_3(P_0)_3(\rho_0)_3$$

$$= -\mu_3(a_3)u_3 - (f_3)_3(P_3) + (f_2)_3(Q_3))u_3 - (\rho_0)_3 (f_3)_3(P_3) - (f_3)_3(P_0)_3$$

$$+ (f_2)_3(Q_3)), \qquad (15)$$

Where $(\rho_0^3)_3(a_3)$ denotes $d(\rho_0)_3(a_3)/da_3$

Expansion $(f_3)_3(P_3)$ and $(f_2)_3(Q_3)$ at $(P_0)_3$ and 0 yield respectively

$$(f_3)_3(P_3) = (f_3)_3((P_0)_3 + N_3)$$

$$= (f_3)_3(P_0)_3 + (f_3^3)_3(P_0)_3 N_3 + o_3(N_3), \qquad (16)$$

$$(f_2)_3(Q_3)$$

$$= (f_2^3)_3(0)Q + o_3(Q_3) \qquad (17)$$

Write $f_3(P_0) = f_{30}, f_3'(P_0) = f_{30}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_3)_3 + (u_a)_3 \approx -(\mu_3(a_3) + (f_{30})_3)u - (f_{30}^3)_3(\rho_0)_3(N_3) - (f_{20}^3)_3(\rho_0)_3(Q_3)$$

$$u_3(t, 0) = \rho_3(t, 0)(\rho_0)_3(0)$$

$$= h_3(Q_3)\int_0^{\infty} \beta_3(a_3)(u_3(t,a) + (\rho_0)_3(a_3)da_3 - (\rho_0)_3\,(0)$$

$$= h_3(Q_3)\int_0^{\infty} \beta_3(a_3)(u_3(t,a_3)da_3 + (h_3(Q_3) - 3)(\rho_0)_3\,(0) \qquad (18)$$

Expanding

$$h_3(Q_3) = 1 + \ (h^3)_3(0)Q_3 + 0(Q_3) \equiv 1 + h_0^{'}Q_3 + 0Q_3$$

And substituting it into (38), we arrive at

$$u_3(t,0) \approx \int_0^{\infty} \beta_3\ (a_3)u_3(t,a_3)da_3 + (h_{a_3}^{'})(\rho_0)_3Q_3$$

The linearization is the following :

$$[u_3] + [u_3] = -(\mu_3(a_3) + [t_{30}]_3u_3 - [\alpha_3]_3(a_3)N_3 - [\alpha_2]_3(a_3)Q_3$$

$$u_3(t,0) = \int_0^{\infty} \beta_3(a_3)u_3(t,a_3)da_3 + [\xi_3]Q_3 \qquad (19)$$

$$\frac{dQ_3}{dt} = -b_3Q_3$$

where

$$[\alpha_3]_3(a_3) = [t_3^3a]_3[p,o]_3(a_3)d2(a_3) = [f_{20}^3]_3(\rho\ o\ )_3(a_3)\xi = [\mathrm{h\ o}^3]_3[\rho\mathrm{o}]_3(\mathrm{o}) \text{ and } \mathrm{b}_3 = \mathrm{b}_3 - g_3(P_3)$$

If $g_3(P_0) > b$, clearly, $(\rho_0)_3(a_3), 0)$ is unstable and the predator population is persistent. Assume $g_3(P_0)_3 < b$.

Then, $Q_3(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_3 + u_a = -(\mu_3(a_3) + (f_{30})_3)u_3 - (\alpha_3)_3(a_3)N_3$$

$$= -(\mu_3(a_3) + f_{30})u_3 - \alpha_3(a_3)\int_0^\infty u_3(t, a_3)da_3$$

$$u_3(t, 0) = \int_0^\infty \beta_3\,(a_3)u_3(t, a_3)da_3$$

Using the method developed by Cushing [2, 3], we suppose

$$u_3(t, a_3) = \tilde{u}_3(a_3)^{e^{\gamma(t-a_3)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_3}\,da_3$. Then,

A solution of the first differential equation for $\tilde{u}(a_3)$ gives,

$$\tilde{u}_3(a_3) = e^{-\int_0^a \mu_3(\tau_3)d\tau - (f_{30})_3 a}\left(\tilde{u}_3(0) - w\int_0^a e^{\int_0^\tau \mu_3(s_3)ds + (f_{30})}\tau_3\,(\alpha_3)_3(\tau_3)e^{\gamma_3\tau}d\tau\right).$$

Substituting the expression (13) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\right)\tilde{u}_3(0)$$

$$+w\int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\int_0^a (\alpha_3)_3(\tau_3)e^{-\gamma_3\tau_3}e^{(f_{30})_3\tau_3 + \int_0^\tau \mu_3(s_3)ds}d\tau_3\,da_3$$

$$= 0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(\tau_3)d\tau} da\right)(\tilde{u}_3(0))_3$$

$$+\left(-3-\int_0^\infty e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(\tau_3)d\tau}\int_0^a e^{(f_{30})_3\tau_3+\int_0^\tau \mu_3(s_3)ds+y\tau}\alpha_3(\tau)d\tau\, da\right)w=0. \quad (23)$$

Define $Y(\gamma,a)=e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(t_3)d\tau}$. Then

$$\left(1-\int_0^\infty \beta_3\,(a_3)Y_3(\gamma_3,a_3)da_3\right)u_3(0)$$

$$+w\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)\int_0^a(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\, da_3=0, \quad (24)$$

$$\left(\int_0^\infty Y_3(\gamma_3,a_3)da_3\right)\tilde{u}_3(0)$$

$$+\left(-1-\int_0^\infty Y_3(\gamma_3,a_3)\int_0^a(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\, da_3\right)w=0.$$

Equqtion (20) has a solution of the form of $u_3(t,a_3)=u_3(a_3)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_3(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_3(\gamma_3,a_3)=\int_0^a(\rho_0)_3(\tau_3)Y^{-3}\,(\gamma_3,a_3)d\tau_3$. We obtain the characteristic equation

$$f_{30_3}^l\left(\int_0^\infty Y_3(\gamma_3,a_3)da_3\right)\left(\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$+\left(1-\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)da\right)\left(1+(f_{30}^l)_3\int_0^\infty Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_3(a_3)=(\rho_0)_3(0)e^{-\int_0^a \mu_3(\tau_3)d\tau_3-(f_{30})_3a_3}$$

$$I_3(\gamma_3,a_3)=\int_0^a (\rho_0)_3(\tau_3)e^{\gamma_3\tau_3+(f_{30})_3\tau+\int_0^\tau \mu_3(s_3)ds_3}d\tau$$

$$=\frac{(\rho_0)_3(0)}{\gamma_3}(e^{\gamma_3a_3}-1)$$

Let

$$E_3(a_3)=e^{-(f_{30})_3a_3-\int_0^a \mu_3(\tau_3)d\tau_3}$$

Then $[Y_3]_3[\gamma_3]_3,a_3)=E_3(a_3)e^{-\gamma_3a_3}$. by a simple calculation,

$$\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)I(\gamma_3,a_3)da=\frac{[\rho_0]_3(0)}{\gamma}\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3a_3}\ da\right);$$

$$\int_0^\infty Y_3(\gamma_3,a_3)I(\gamma_3,a_3)da=\frac{[\rho_0]_3(0)}{\gamma_3}\left(\int_0^\infty E_3(a_3)e^{-\gamma_3a_3}\ da_3\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_3(0)}{\gamma_3}[f_{30}^3]\int_0^\infty E_3(a_3)e^{-\gamma_3a_3}da_3\left(\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3a_3}da_3\right)$$

$$+\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right).$$

$$\left(1+\frac{(\rho_0)_3(0)}{\gamma_3}(f_{30}^{'})_3\left(\int_0^\infty E_3(a_3)da_3-\int_0^\infty E_3(a_3)e^{-\gamma_3 a_3}da_3\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right)(1+\frac{(f_{30}^{'})_3}{\gamma_3}(\rho_0)_3(0)\int_0^\infty E_3(a_3)da_3)=0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_3(a_3)Y_3(\gamma_3 a_3)da_3 = 1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3 \neq 0,$$

Since $\int_0^\infty \beta_3(a_3)Y_3(\gamma_3 a_3)\int_0^a \alpha_3(\tau_3)Y^{-3}(\gamma_3,\tau_3)d\tau_3 da_3 > 0$. Therefore, the characteristic solution is

$$\gamma_3 = -(f_{30}^{'})_3(\rho_0)_3\int_0^\infty E_3(a_3)da_3.$$

Since $f_3'(P_0)_3 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_3), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_3), 0)$ which is related to the extinction of debit summation types

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_3 + (\mu_3(a_3) + [f_3]_3(P_3^*) + [f_2]_3(Q_3^*))P_3^* = 0, \quad (26)$$

$$p_3^*(0) = h_3(Q_3^*)\int_0^\infty \beta_3(a_3)p_3^*(a_J)p_3^*(a_3)da_3 \quad (27)$$

$$g_3(P_3^*) - b_3 = 0, \quad (28)$$

Where $P_3^* = \int_0^\infty p_3^*(a_3)\, da_3$

Equation (28) yields $P^* = g^{-3}(b)$,

From (26),

$$p_3^*(a) = p_3^*(0)e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}$$

Substituting into (27), we need to solve

$$\left(1 - h_3(Q_3^*)\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}da_3\right)p_3^*(0) = 0$$

For $Q_3^*$

Fix $p_3^*(0)$ and $P_3^*$ Let

$$H_3(z_3) = 1 - h_3(z_3)\int_0^{\infty} \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3} da_{,3}$$

And assume that

$$[k_3]_3 \int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(g_3^{-3}(b_3))a_3 - [k_2]_3 a_3}\, da_3 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_3(z_3)] > 0.$$

Clearly,

$$\frac{dH_3}{dz_3} = h_3(z)([f_2^{'}]_3(z_3))\int_0^{\infty} a_3\beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - f_3(P_3^*)a_3 - f_2(z_3)a_3} da_{\square}$$

$-h^{'}(z_3)\int_0^{\infty} \beta_3(a_3)\, e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3}\ da_3$

By the assumptions on $\left[f_2^3\right]_3 and\ [h^3]^3 dH_3/dz > 0$. Hence, if

$H_3(0) = 1 - \int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3} da_3 < 0,$ (29)

There exists a unique positive solution $Q^*$

Now, we claim that $H_3(0) < 0$ follows if $g_3^{-}(b_3) < [f_3^{-3}]_3(r_3)$, where $r_3$ is the solution of (12).

Indeed, since

$$R_3(s_3) = \int_0^{\infty} \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 e^{-y\frac{a}{3}}} da_3$$

Is a decreasing function of $s_3$ when $P_3^* = [g^{-3}]_3(b_3) < [f_3^{-3}]_3(r_3)$, that is $[f_3]_3(P_3^*) < r$, $R\big(f_3(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_3(g^{-3}(b))) > 1, i.e.,$

$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3\left(g_3^{-3}(b)\right)a_3} > 1.$ (H4)

Then the system (4)-(6) has a unique positive equilibrium $(P_3^*(a_3), Q_3^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_3^*(a_3), Q_3^*)$ is more complicated. We again consider perturbations $m_3(t_3 a_3) = p_3(t_3, a_3) - p_3^*(a_3)$ and $W_3(t) = Q_3(t) - Q^*$.Then

$M_3(t) \equiv \int_0^{\infty} m_3(t, a_3) da_3 = P_3(t) - P_3^*,$

$$[m_3]_3 + [m_a]_3 = [p_3]_3 + [p_a]_3 - \left[p^*\right]_3(a_3) \approx -(\mu_3(a_3) + [f_{3*}]_3 + [f_{2*}]_3)m_3$$

$$-[f_{3*}^3]_3 p_3^*(a_3)M - [f_{2*}^3]_3 p_3^*(\mathrm{a})W_3,$$

$$m_3(t,0)p_3(t,0) - p_3^*(0) \approx h_* \int_0^{\infty} \beta_3(a_3) m_3(t, a_3)\, da + \frac{[h_*^3]_3}{[h_*]_3}\, p_3^*(0)W_3$$

$$\frac{dW_3}{dt} = \frac{dQ_3}{dt} \approx [g_*^3]_3 Q_3^* M_{,3}$$

Where $[f_{3*}]_3$ denoting $[f_3(P^*)]_3$, $[f_{3*}]_3$ denoting $[f_2(P^*)]_3$, $[f'_{3*}]_3$ denoting $[f'_3(P^*)]_3[f'_{2*}]_3$ denoting $[f_2]_3$ $(Q_3^*,\quad h_3^*$ denoting $h'_*$ denoting $h'_3$ $(Q_3^*)$, and $[g'_*]_3$

Denoting $[g']_3(P_3^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_3]_3 + [m_a]_3 = -(\mu_3(a_3) + [f_{3*}]_3 + [f_{2*}]_3)m_3 - [f'_{3*}]_3[p_3^*](a_3)M_3 - [f'_{2*}]_3 p_3^*(a_3)W_3$$

$$m_3(t,0) = [h_*]_3 \int_0^\infty \beta_3(a_3)m_3(t,a_3)da_3 + \frac{[h_*^3]_3}{[h^*]_3}[p^*]_3(0)\,W_{,3} \qquad (30)$$

$$\frac{dW_3}{dt} = [g_*^3]_3[Q_3^*]M_3$$

Integrating $[m_3]_3(t,a_3)$ with respect to a from 0 to ∞, we have

$$\frac{dM_3(t)}{dt} = \int_0^\infty [mt]_3(t,a_3)da_3$$

$$= -\int_0^\infty [m_a]_3\,(t,a_3)da_3 - \int_0^\infty \mu(a_3)m(t,a_3)da_3$$

$$-\left([f_{3*}]_3 + [f_{2*}]_3 + [f'_{3*}]_3\int_0^\infty p_3^*(a_3)da_3\right)M_3 - [f_{2*}^3]_3\int_0^\infty p_3^*(a_3)da_3 W_3$$

$$= m_3(t,0) - \int_0^\infty \mu_3(a_3)\,m_3(t,a_3)da_3$$

$$-\left([f_{3*}]_3 + [f_{2*}]_3 + [f'_{3*}]_3 P_3^*\right)M_3 - [f'_{2*}]_3 P_3^* W_{,3}$$

From (30), we arrive at the equations

$$\frac{dW_3(t)}{dt} = \int_0^{\infty} ([h_*]_3\beta_3(a_3) - \mu_3(a_3))m_3(t,a_3)da_3$$

$$-\left([f_{3*}] + [f_{2*}]_3 + \left[f_{3*}^{'}\right]_3 P_3^*\right) M_3 + \left(\frac{[h_*^{'}]_3}{[h_*]^3}\, P_3^*(0) - f_{2*}^{'} P_3^*\right) W_3,$$

(33)

$$\frac{dW_3(t)}{dt}\left[g^{'}{}_*\right]_3 Q_3^*\, M_{.3}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{"*}]_3\beta_3(a_3) - \mu_3(a_3))m_3(t,a_3)da_3 = ([h_*]_3\beta_3(\theta_3) - \mu_3(\theta_3))M_{,3}$$

Where $0< \theta_3 = \theta_3(t) < \infty$.

And if $a_3^2\,[a_3^2]_3<4\,[a_2]_3<4\,a_2 e$,

$\mathrm{Re}\,\lambda_3\left[\lambda_3\right]_3 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_3 = -\frac{3}{2}a_3[a2\ ]_3 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_*\,[h_*]_3\beta_3(a_3) - \mu_3(a_3)\right|$. Then$|n_3(t)|$ ⊠⊿∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_3^2 <4\,a_2 e$ and $[a_3] > 2\Delta$,or $[a_3^2]$ ⊠$[a_2]e$ and $(a_3 - \sqrt{(a_3^2)_3 - 4(a_2)_3 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+3}\right)$ and $\beta_3(a_3) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3) + d\tau_3 da} = \frac{3}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_3 + [\mu_0]$ 0.5. Assume $[f_3](x_3) = K_3 x_3$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_3(x_3) = [0x_3]/(3 + x_3)$; the functional responses $[[f_2](x) = [0]\, x_3/(1 + x_3)$ and $h_3(x_3)=(1+\omega_3 x_3)/(1 + x_3)$ (See Holling [7] and Rashevsky [13])

Equation (34),now has the form of

$$\frac{m_3 + [\mu_0] + [K][P_0])}{[m_3]_3 + (\mu_0)_3 + [K][P_0])_3^3\ [K]_3 [P_0]_3^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(3 + P_0) < b_3$ i e $[P_0] < b/(\delta_3 - b_3)$ then the equilibrium $(P_0)$ $[P_0]_3(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_3$

$[f_{3*}] = K_3\ P^*\ [f_{2*}] = \frac{\sigma Q_3^*}{3+Q_3^*} \quad [h_*] = \frac{3+\omega_3 Q_3^*}{3+Q_3^*}$

$[f_{3*}] = K_3\ [f_{2*}^3] = \frac{\sigma_3}{(3+Q_3^*)^2}\ [g_*^3] = \frac{\sigma_3}{(3+P_3^*)^2} \quad [h_*^3] = \frac{\omega_3 - 3}{(3+Q_3^*)^2}$

In order to have (H3) and (H4) hold, $m, \mu_0, K, b, \delta$, and $\omega$ need to satisfy

$$\frac{\omega_3}{(m+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma)^2}+\frac{2\omega_3}{([m_3]_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma_3)^3}<1$$

And

$$\frac{1}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^2}+\frac{2\omega_3}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}+(m_3+[\mu_0]_3+\frac{K_3b_3}{\delta_3-b_3}+\sigma_3)\,Q_3^*)^3$$

$$=m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}(1+Q^*)(1+\omega_3Q_3^*$$

Now

$$\left[a_3^2\right]-4\left[a_2\right]e=\left(\frac{2K_3b_3}{\delta_3-b_3}\right)^2+\frac{4K_3b_3\sigma_3\,Q_3^*}{(\delta_3-b_3)(1+Q_3^*)}+\frac{\sigma_3^2Q_3^{*2}}{(1+Q_3^*)^2}$$

$$\frac{4\,b_3\sigma_3\,(\delta_3-b)Q^*}{\delta_3(1+Q_3^*)^2}-\frac{4\,(\delta_3-b_3)^2(1-\omega_3)Q_3^*}{\delta_3(1+\omega_3Q_3^*)(1+Q_3^*)}P*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_3\,ae_3^{mu}\;\;[\mu_0]_3-\frac{a_3}{a_3+3}$$

Hence if $[a_3^2 > 4(a_2)_3 e$ and

$$[a_3] = \frac{2K_3 b_3}{\delta_3 - b_3} + \frac{\sigma_3 Q_3^*}{1 + Q_3^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_3]_3 = [f_{3*}]_3 + [f_{2*}]_3 + [f_{3*}] P_3^*$$

$$[a_2]_3 = \left[f_{2*}^3\right] P_3^* - \frac{h_*^3}{h_*} P^*(0)$$

$$e = [g_*^3]_3 Q_3^*$$

$$\eta_3(t) = \left[h_*\right]_3 \beta_3 \left(\theta_3(t)\right) - \mu_3(\theta(t)),$$

The system(33)can be written as the following system:

$$[M(t)]\left[a_3\right] M_3 - [a_2]_3 W_3 - [n(t)] M_3$$

$$w_3(t) = eM_3$$

Write

$$A = \begin{pmatrix} -[a_3]_3 - [a_2]_3 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_3 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_3 \\ W_3 \end{pmatrix}$$

The solution of (32)is

$$X_3(t) = e^{A,t}[X]_3(0) + \int_0 e^{At}$$

**This equation is useful in following type of problems . For actual model kindly see various papersThere is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant Λ, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive Λ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the Λ>0 generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible.**

**Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of (e) GW signals from BBHs with component spins $\chi_{1,2} \in [-0.99, 0.99]$ aligned (e&eb) with the orbital angular momentum, component masses $m_{1,2} \in [2,48] M_\odot$, and (e&eb) total mass $M_{total} \leq 50\ M_\odot$. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins >+0.6 and component masses $\gtrsim 20\ M_\odot$. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201 STOP Advanced LIGO may be the first experiment to (e) detect gravitational waves. Through superradiance of stellar black holes, it may also be (=) the first experiment to discover the QCD axion with (e&eb) decay constant above the GUT scale. When an axion's Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e) the black hole, forming a "gravitational atom." Through the superradiance process, the number of axions occupying (e) the bound levels grows exponentially, extracting (e) energy and angular momentum from the black hole. Axions**

**transitioning between levels of the gravitational atom and (e&eb) axions annihilating to (e) gravitons can produce (eb) observable gravitational wave signals. The signals are long-lasting, monochromatic, and can be distinguished from (e) ordinary astrophysical sources. Authors estimate up to O (1) transition events at aLIGO for (e) an axion between 10^-11 and (e&eb) 10^-10 eV and up to 10^4 annihilation events for (e) an axion between 10^-13 and 10^-11 eV. In the event of a null search, aLIGO can constrain (e) the axion mass for (e) a range of rapidly spinning black hole formation rates. Axion annihilations are also promising for (e) much lighter masses at future lower-frequency gravitational wave observatories; the rates have (e) large uncertainties, dominated by (e) supermassive black hole spin distributions. Projections for aLIGO are (=) robust against perturbations from (e) the black hole environment and account for (e) updated exclusion on (e&eb) the QCD axion of 6*10^-13 eV < ma < 2*10^-11 eV suggested by (e) stellar black hole spin measurements. Subjects: High Energy Physics - Phenomenology (hep-ph); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 91, 084011 (2015) DOI: 10.1103/PhysRevD.91.084011Cite as:arXiv:1411.2263 [hep-ph] (or arXiv:1411.2263v3 [hep-ph] for this version) Discovering the QCD Axion with Black Holes and Gravitational Waves Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang STOP Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric**

**charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by (e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave**

**signal. The signal sweeps upwards in frequency from 35 to 250 Hz with (e&eb) a peak gravitational-wave strain of (e) 1.0×10−21. It matches the waveform predicted by general relativity for (e) the inspiral and merger of a pair of black holes and the ringdown of the resulting (eb) single black hole. The signal was observed with a matched-filter signal-to-noise ratio of (e) 24 and a false alarm rate estimated to be (=) less than 1 event per 203 000 years, equivalent to (=) a significance greater than 5.1σ. The source lies at (eb) a luminosity distance of 410+160−180 Mpc corresponding to (e&eb) a redshift z=0.09+0.03−0.04. In the source frame, the initial black hole masses are (=) 36+5−4M⊙ and 29+4−4M⊙, and the final black hole mass is (=) 62+4−4M⊙, with (e&eb) 3.0+0.5−0.5M⊙c2 radiated in (eb) gravitational waves. All uncertainties define (eb) 90% credible intervals. These observations demonstrate (eb) the existence of binary stellar-mass black hole systems. Observation of Gravitational Waves from a Binary Black Hole Merger B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 061102 – Published 11 February 2016 This is the first direct detection of gravitational waves and (e&eb) the first observation of a binary black hole merger. Received 21 January 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.061102 STOP This is the second in a series of three papers on (e&eb) systematic analysis of rank 1 Coulomb branch geometries of (e) four dimensional N=2 SCFTs. In the first paper authors developed a strategy for (e) classifying physical rank-1 CB geometries of N=2 SCFTs. They show how to carry out this strategy computationally to construct (eb) the Seiberg-Witten curves and (e&eb) one-forms for all the rank-1 SCFTs. Explicit expressions are given for all cases, with the exception of the Nf=4 SU (2) gauge theory and (e&eb) the En SCFTs which were previously constructed. Classification includes all known rank-1 theories plus (+) a new one with (e&eb) an abelian flavor group, plus (+) nine additional theories whose existence is more speculative. Four of those, reported in first paper, depend on (=) the assumption of new frozen rank-1 SCFTs. Here authors also show that the assumption of the existence of certain rank-0 N=2 SCFTs leads to (eb) five additional consistent**

**rank-1 CB geometries. Subjects: High Energy Physics - Theory (hep-th) Cite as:arXiv: 1601.00011 [hep-th] (or arXiv: 1601.00011v2 [hep-th] for this version Geometric constraints on the space of N=2 SCFTs II: Construction of special Kähler geometries and RG flows Philip C. Argyres, Matteo Lotito, Yongchao Lü, Mario Martone (Submitted on 31 Dec 2015 (v1), last revised 6 Jan 2016 (this version, v2)) STOP Authors propose a model of Dark Energy in which the field currently dominating (e) the energy density of the universe is (=) an "axion field" linearly coupled to (e&eb) the Pontryagin density, tr (FΛF), (i.e., the exterior derivative of (e) the Chern-Simons form) of (e) a massive gauge field. Authors assume that the axion has (e) self-interactions corresponding to (e&eb) a non-trivial (exponential) potential. Authors argue that a non-vanishing magnetic helicity of (e) the gauge field triggers (eb) slow-rolling of the axion at field values far below (eb) the Planck scale. Proposal leads to (eb) a "Tracking Dark Energy Scenario" in which the contribution of the axion energy density to (e) the total energy density is (=) constant (and small) during (e&eb) the early radiation phase, until (e) a secular growth term proportional to (e) the Pontryagin density of (e) the gauge field becomes (=) dominant. The initially small contribution of (e) the axion field to (e) the total energy density is related to the observed small baryon-to-entropy ratio. Some speculations concerning the nature of the gauge field are offered.**

**Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Cite as:arXiv: 1601.00057 [hep-th] (or arXiv:1601.00057v2 [hep-th] for this version) Tracking Dark Energy from Axion-Gauge Field Couplings Stephon Alexander, Robert Brandenberger, Juerg Froehlich (Submitted on 1 Jan 2016 (v1), last revised 23 Jan 2016 (this version, v2)) STOP Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a**

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**associated entropic measure. Subjects: High Energy Physics - Theory (hep-th) DOI:10.1016/j.physletb.2016.02.038 Cite as: arXiv: 1601.00076 [hep-th] (or arXiv: 1601.00076v2 [hep-th] for this version) Bounds on topological Abelian string-vortex and string-cigar from information-entropic measure R. A. C. Correa, D. M. Dantas, C. A. S. Almeida, Roldao da Rocha (Submitted on 1 Jan 2016 (v1), last revised 17 Feb 2016 (this version, v2)) STOP Sum of the Holevo quantity (that bounds the capacity of quantum channels to transmit (e&eb) classical information about an observable) and (e&eb) the quantum discord (a measure of the quantumness of (e) correlations of that observable) yields (eb) an observable-independent total given by (e) the quantum mutual information This split naturally delineates ,dilates and disseminates information about (e&eb) quantum systems accessible to (e) observers -- information that is redundantly transmitted by (e&eb) the environment -- while showing that it is (=) maximized for (e) the quasi-classical pointer observable. Other observables are accessible only via (e&eb) correlations with the pointer observable .Further authors prove an anti-symmetry property relating accessible information and (e&eb) discord. It shows that information becomes (=) objective -- accessible to (e) many observers -- only as quantum information is relegated to (e) correlations with (e&eb) the global environment, (eb) and, therefore, locally inaccessible. The resulting complementarity explains (eb) why, in a quantum Universe, we perceive objective classical reality while (e) flagrantly quantum superpositions are out of reach. Subjects: Quantum Physics (quant-ph) Journal reference: Scientific Reports 3, 1729 (2013) Report number: LAUR-13-26386 Cite as:arXiv: 1303.4659 [quant-ph] (or arXiv: 1303.4659v2 [quant-ph] for this version) Complementarity of quantum discord and classically accessible information Michael Zwolak, Wojciech H. Zurek STOP Niels Bohr proposed that the outcome of the measurement becomes (=) objective and real, and, hence, classical, when (e) its results can be communicated by (e) classical means. In this work authors revisit Bohr's postulate using (e) modern tools from (e) the quantum information theory. They find a**

**full confirmation of Bohr's idea: if a measurement device is in a nonclassical state, (eb) the measurement results cannot be communicated perfectly by classical means. In this case some part of information in the measurement apparatus is lost (e) in the process of communication: the amount of this lost information turns out to be (=) the quantum discord. The information loss occurs even when (e) the apparatus is not entangled with the system of interest. The tools presented in this work allow (eb) to generalize Bohr's postulate: authors show that for pure system-apparatus states (eb) quantum communication does not provide any advantage when (e) measurement results are communicated to (e) more than one recipient. They further demonstrate the superiority of quantum communication to (e&eb) two recipients on (e&eb) a mixed system-apparatus state and show that (eb) this effect is fundamentally different from (e) quantum state cloning. Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. Lett. 111, 040401 (2013) DOI:10.1103/PhysRevLett.111.040401 Report number: LA-UR-12-25351 Cite as: arXiv: 1304.2914 [quant-ph] (or arXiv: 1304.2914v2 [quant-ph] for this version) Quantum discord cannot be shared Alexander Streltsov, Wojciech H. Zurek STOP Wojciech H. Zurek, Don N. Page investigate thermodynamic equilibrium of (e) a self-gravitating perfect fluid in (eb) a spherically symmetric system containing (e) a black hole of mass M by means of (e) the Tolman-Oppenheimer-Volkoff (TOV) equation. At r >> 2M its solutions describe (eb) a black-body radiation atmosphere with (e&eb) the Hawking temperature T_BH~1/(8 \pi M) that is increasingly blueshifted as r approaches (eb) 2M. However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments:**

**Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by**

**(e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a firewall-like "Planck cocoon" with entropy close to (e) black hole entropy**

# Prescientiаl Dynamics And Contential Staticity Of Accentuation Attrition Models

## Module Four

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equailibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-445) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Some of the concepts to which model is applicable is as follows:**

**Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by (e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave signal. The signal sweeps upwards in frequency from 35 to 250 Hz with (e&eb) a peak gravitational-wave strain of (e) 1.0×10−21. It matches the waveform predicted by general relativity for (e) the inspiral and merger of a pair of black holes and the ringdown of the resulting (eb) single black hole. The signal was observed with a**

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They further demonstrate the superiority of quantum communication to (e&eb) two recipients on (e&eb) a mixed system-apparatus state and show that (eb) this effect is fundamentally different from (e) quantum state cloning. Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. Lett. 111, 040401 (2013) DOI:10.1103/PhysRevLett.111.040401 Report number: LA-UR-12-25351 Cite as: arXiv: 1304.2914 [quant-ph] (or arXiv: 1304.2914v2 [quant-ph] for this version) Quantum discord cannot be shared Alexander Streltsov, Wojciech H. Zurek STOP Wojciech H. Zurek, Don N. Page investigate thermodynamic equilibrium of (e) a self-gravitating perfect fluid in (eb) a spherically symmetric system containing (e) a black hole of mass M by means of (e) the Tolman-Oppenheimer-Volkoff (TOV) equation. At r >> 2M its solutions describe (eb) a black-body radiation atmosphere with (e&eb) the Hawking temperature T_BH~1/(8 \pi M) that is increasingly blueshifted as r approaches (eb) 2M. However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments: Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by (e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a**

**firewall-like "Planck cocoon" with entropy close to (e) black hole entropy. Authors reproduce paper with an updated title and abstract. For a later study with similar results see G. 't Hooft, gr-qc/9706058 Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Physical Review D, Volume 29, Number 4, pp. 628 - 631 (1984) DOI:10.1103/PhysRevD.29.628 Cite as:arXiv:1511.07051 [gr-qc] (or arXiv:1511.07051v1 [gr-qc] for this version) Firewalls, black-hole thermodynamics, and singular solutions of the Tolman-Oppenheimer-Volkoff equation Wojciech H. Zurek, Don N. Page (Submitted on 22 Nov 2015) STOP Frederic Magniez, Ashwin Nayak consider the problem of testing the commutativity of (e) a black-box group specified by its k generators. The complexity (in terms of k) of (e) this problem was first considered by Pak, who gave (eb) a randomized algorithm involving O(k) group operations. They construct a quite optimal quantum algorithm for this problem whose complexity is in O~(k2/3)O~(k2/3). The algorithm uses and highlights the power of the quantization method of Szegedy. For the lower bound of Ω(k2/3)Ω(k2/3), authors give a reduction from (e) a special case of Element Distinctness to our problem. Along the way, they prove the optimality of the algorithm of (e) Pak for the randomized model. Algorithmica July 2007, Volume 48, Issue 3, pp 221-232 First online: 25 June 2007 Quantum Complexity of Testing Group Commutativity Frederic Magniez, Ashwin Nayak STOP Pawel Wocjan, Jon Yard analyze relationships between quantum computation and (e&eb) a family of generalizations of the Jones polynomial. Extending recent work by Aharonov et al., authors give efficient quantum circuits for (e) implementing the unitary Jones-Wenzl representations of the braid group. We use these to provide new quantum algorithms for (e) approximately evaluating a family of specializations of the HOMFLYPT two-variable polynomial of trace closures of braids. They also give (eb) algorithms for approximating the Jones polynomial of a general class of closures of (e) braids at roots of unity. Next authors provide (eb) a self-contained proof of a result of (e) Freedman et al. that any quantum computation can be replaced by (e&eb) an additive approximation of the Jones polynomial, evaluated at (eb) almost any primitive root of unity. Proof encodes two-qubit unitaries into (e&eb) the rectangular representation of (e) the eight-strand braid group. Authors then give QCMA-complete and (e&eb) PSPACE-complete problems which are based on braids. They conclude with direct proofs that evaluating the Jones polynomial of the plat closure at most primitive roots of unity is a #P-hard problem, while learning its most significant bit is PP-hard, circumventing (e) the usual route through the Tutte polynomial and graph coloring. Subjects: Quantum Physics (quant-ph); Geometric Topology (math.GT) Cite as: arXiv:quant-ph/0603069 (or arXiv:quant-ph/0603069v3 for this version) The Jones polynomial: quantum algorithms and applications in quantum complexity theory Pawel Wocjan, Jon Yard STOP Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan address the task of switching off the Hamiltonian of (e) a system by removing (e) all internal and system-environment couplings. Authors propose dynamical decoupling schemes, which use (e) only bounded-strength controls, for quantum many-body systems with local system Hamiltonians and local environmental couplings. To do so, they introduce (eb) the combinatorial concept of (e) balanced-cycle orthogonal arrays (BOAs) and show how to construct them from (e) classical error-correcting codes. The derived decoupling schemes may be useful as (=) a primitive for more complex schemes, e.g., for Hamiltonian simulation. For the case of n qubits and a 2-local Hamiltonian, the length of the resulting decoupling (e&eb) scheme scales as O (nlogn), improving over (e&eb) the previously best-known schemes that scaled quadratically with (e&eb) n. More generally, using balanced-cycle orthogonal arrays constructed from families of BCH codes, we show that bounded-strength decoupling for any ℓ-local Hamiltonian, where ℓ≥2, can be achieved using decoupling schemes of length at most O(nℓ−1logn). Subjects:Quantum Physics (quant-ph); Emerging Technologies (cs.ET) Journal reference: IEEE Trans. Inf. Theory, 62(5): 2881-2894, 2016 DOI: 10.1109/TIT.2016.2535183 Report number: MIT-CTP /4705 Cite as: arXiv: 1509.00408 [quant-ph] (or arXiv: 1509.00408v2 [quant-ph] for this version) Improved bounded-strength decoupling schemes for local Hamiltonians Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan STOP Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola propose dynamical control schemes for Hamiltonian simulation in many-body quantum systems that avoid instantaneous control operations and rely solely on (e) realistic bounded-strength control Hamiltonians. Each simulation protocol consists of (e) periodic repetitions of a basic control block, constructed as (=) a suitable modification of an "Eulerian decoupling cycle," that would otherwise implement (eb) a trivial (zero) target Hamiltonian. For an open quantum system coupled to (e) an uncontrollable environment, approach may be employed to (e) engineer an effective evolution that simulates (e&eb) a target Hamiltonian on the system, while suppressing (e) unwanted decoherence to the leading order. Authors present illustrative applications to (e) both closed- and open-system simulation settings, with (e&eb)**

**emphasis on simulation of non-local (two-body) Hamiltonians using (e) only local (one-body) controls. In particular, authors provide (eb) simulation schemes applicable to (e&eb) Heisenberg-coupled spin chains exposed to (e) general linear decoherence, and show how to simulate (e&eb) Kitaev's honeycomb lattice Hamiltonian starting from Ising-coupled qubits, as (=) potentially relevant to (e&eb) the dynamical generation of a topologically protected (e&eb) quantum memory. Additional implications for quantum information processing are discussed. Subjects: Quantum Physics (quant-ph) Journal reference: New J. Phys. 16, 045021 (2014) DOI: 10.1088/1367-2630/16/4/045021 Cite as:arXiv: 1310.4153 [quant-ph]**

**(or arXiv: 1310.4153v1 [quant-ph] for this version) Hamiltonian quantum simulation with bounded-strength controls Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola STOP Felix Fontein, Pawel Wocjan present a quantum algorithm for computing the period lattice of (e) infrastructures of fixed dimension. The algorithm applies to (e&eb) infrastructures that satisfy certain conditions. The latter are always fulfilled for (e) infrastructures obtained from (e) global fields, i.e., algebraic number fields and function fields with (e&eb) finite constant fields. The first main contributions is an exponentially better method for (e) sampling approximations of vectors of the dual lattice of (e) the period lattice than the methods outlined in the works of (e) Hallgren and Schmidt and Vollmer. This new method improves (eb) the success probability by (e) a factor of at least 2^ {n^2-1} where (e) n is the dimension. The second main contribution is a rigorous and complete proof that the running time of the algorithm is (=) polynomial in the logarithm of (e) the determinant of the period lattice and (e&eb) exponential in n. The third contribution is the determination of (e) an explicit lower bound on (e&eb) the success probability of our algorithm which greatly improves on (e&eb) the bounds given in the above works. The exponential scaling seems inevitable because (e) the best currently known methods for (e) carrying out fundamental arithmetic operations in infrastructures obtained from (e) algebraic number fields take exponential time. In contrast, the problem of computing (e&eb) the period lattice of infrastructures arising from function fields can be solved without (e) the exponential dependence on (e&eb) the dimension n since this problem reduces efficiently to the abelian hidden subgroup problem. This is also true for (e) other important computational problems in (eb) algebraic geometry. The running time of the best classical algorithms for (e) infrastructures arising from (e) global fields increases subexponentially with (e&eb) the determinant of the period lattice. Subjects: Quantum Physics (quant-ph) Cite as: arXiv: 1111.1348 [quant-ph]**

**(or arXiv: 1111.1348v2 [quant-ph] for this version) Quantum Algorithm for Computing the Period Lattice of an Infrastructure Felix Fontein, Pawel Wocjan STOP Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan advocate a new approach of (e) addressing hidden structure problems and (e&eb) finding efficient quantum algorithms. Authors introduce and investigate the Hidden Symmetry Subgroup Problem (HSSP), which is a generalization of the well-studied Hidden Subgroup Problem (HSP). Given a group acting on (e&eb) a set and an oracle whose level sets define (eb) a partition of the set; the task is to recover (e) the subgroup of symmetries of (e) this partition inside the group. The HSSP provides (eb) a unifying framework that, besides the HSP, encompasses (e) a wide range of algebraic oracle problems, including (e) quadratic hidden polynomial problems. While the HSSP can have (e) provably exponential quantum query complexity, authors obtain efficient quantum algorithms for various interesting cases. To achieve this, authors present a general method for (e) reducing the HSSP to the HSP, which works efficiently in several cases related to (e&eb) symmetries of polynomials. The HSSP therefore connects in a rather surprising way certain hidden polynomial problems with (e&eb) the HSP. Using this connection, authors obtain the first efficient quantum algorithm for (e) the hidden polynomial problem for (e) multivariate quadratic polynomials over fields of (e) constant characteristic. Authors also apply the new methods to (e&eb) polynomial function graph problems and (e&eb) present an efficient quantum procedure for (e) constant degree multivariate polynomials over (e&eb) any field. This result improves in several ways the currently known algorithms. Subjects: Quantum Physics (quant-ph) Journal reference: SIAM Journal on Computing 42:5 (2013) DOI: 10.1137/120864416 Cite as: arXiv: 1107.2189 [quant-ph]**

**(or arXiv: 1107.2189v2 [quant-ph] for this version) Hidden Symmetry Subgroup Problems Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan STOP Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan present an efficient general method for realizing a quantum walk operator corresponding to (e&eb) an arbitrary sparse classical random walk. Approach is based on Grover and Rudolph's method for preparing coherent versions of (e) efficiently integrable probability distributions. This method is intended for (e) use in quantum walk algorithms with polynomial speedups, whose complexity is usually measured in terms of (e&eb) how many times we have to apply a step of a quantum walk, compared to (e&eb) the number of necessary classical Markov chain steps. Authors consider a finer notion of complexity including (e) the**

**number of elementary gates it takes to implement each step of the quantum walk with (e&eb) some desired accuracy. The difference in complexity for (e) various implementation approaches is that method scales linearly in the sparsity parameter and (e&eb) poly-logarithmically with (e&eb) the inverse of the desired precision. Best previously known general methods either scale quadratically in (eb) the sparsity parameter, or polynomially in (eb) the inverse precision. Approach is especially relevant for (e) implementing quantum walks corresponding to (e&eb) classical random walks like those used in (eb) the classical algorithms for approximating permanents and sampling from binary contingency tables. In those algorithms, the sparsity parameter grows with (e&eb) the problem size, while (e) maintaining high precision is required. Subjects: Quantum Physics (quant-ph) Journal reference: Quantum Information and Computation, Vol.10, No.5&6, pp0420-0434 (2010) Cite as: arXiv: 0903.3465 [quant-ph] (or arXiv: 0903.3465v3 [quant-ph] for this version) Efficient Circuits for Quantum Walks Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan STOP Preparing the ground state of a system of interacting classical particles is (=) an NP-hard problem. Thus, there is in general no better algorithm to (e) solve this problem than exhaustively going through all N configurations of (e) the system to determine (eb) the one with lowest energy, requiring (e) a running time proportional to N. A quantum computer, if it could be built, could solve (e) this problem in time sqrt (N). Here, authors present a powerful extension of this result to (e) the case of interacting quantum particles, demonstrating (eb)that a quantum computer can prepare the ground state of (e) a quantum system as efficiently as it does for (e) classical systems. Subjects: Quantum Physics (quant-ph) Journal reference: Phys.Rev.Lett.102:130503, 2009 DOI:10.1103/PhysRevLett.102.130503 Cite as: arXiv: 0809.2705 [quant-ph] (or arXiv: 0809.2705v1 [quant-ph] for this version) Preparing ground states of quantum many-body systems on a quantum computer David Poulin, Pawel Wocjan STOP Daniel Nagaj, Pawel Wocjan construct a simple translationally invariant, nearest-neighbor Hamiltonian on (e&eb) a chain of 10-dimensional qudits that makes it possible to (e) realize universal quantum computing without (e) any external control during (e&eb) the computational process. Authors only require the ability to prepare an initial computational basis state which encodes (e) both the quantum circuit and its input. The computational process is (=) then carried out by (e) the autonomous Hamiltonian time evolution. After a time polynomially long in the size of (e) the quantum circuit has passed, the result of (e) the computation is obtained with (e&eb) high probability by (e) measuring a few qudits in the computational basis. This result also implies (eb) that there cannot (e) exist efficient classical simulation methods for (e) generic translationally invariant nearest-neighbor Hamiltonians on (e&eb) qudit chains, unless (e) quantum computers can be efficiently simulated by (e&eb) classical computers (or, put in complexity theoretic terms, unless BPP=BQP). Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. A 78, 032311 (2008) DOI: 10.1103/PhysRevA.78.032311 Cite as: arXiv: 0802.0886 [quant-ph] (or arXiv: 0802.0886v4 [quant-ph] for this version) Hamiltonian Quantum Cellular Automata in 1D Daniel Nagaj, Pawel Wocjan**

## 1. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_4}{dt} = -b_4 Q_4 + g_4(g_4)Q_4$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_4(0) = 0, \quad g_4(P_4) > 0, \ P_4 \ 0, \qquad \lim_{P_4 \to \infty} g_4(P_4) = (K_4)_4 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_4 + (\rho_a)_4 + D_4(a_4, P_4, Q_4)\rho_4 + 0, \qquad 0 < a_4 \ \infty,$$

$$\rho_4(t, 0) = \int_0^\infty n_4(a_4, P_4, Q_4)\rho_4(t_4, a_4)da_4$$

Where $D_4$ is the mortality function and $n_4$ is the fertility function.

Put $_4$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_4}{dt}\right]_4 = x_4\left[\frac{dQ_2}{dt}\right]_4 - x_2\left[\frac{dQ_4}{dt}\right]_4$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_4 - x_4\left[\frac{dQ_2}{dt}\right]_4$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_4 - x_6\left[\frac{dQ_3}{dt}\right]_4$$

Where $x_4, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_4$ and $g_4(P_4)Q_4$

Now when you add the term $(b_4Q_4)$ with $g_4(P_4)(Q_4)$ we are infect see is that $(b_4Q_4)$ is credited incrementally part of $g_4(P_4)(Q_4)$ or instant nearly, when subtracted from $\frac{dQ_4}{dt}$, it means that $g_4(P_4)(Q_4)$ is slowly eaten by $\frac{dQ_4}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_4 = \mu_4(a_4) + (f_4)_4(P_4)_4 + (f_2)_4(Q)_4 \quad (2)$$

$$n_4 = \beta_4(a_4)h_4(Q_4), \quad (3)$$

Where $\mu_4$the age- dependent debit rate whit is $\mu_4(a_4)$ $\forall a > 0$; $(f_4)_4$ is the density- dependent death rate with $f_4(0) = 0$, $(f_4^4)_4(P_4)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_4)_4(P_4) = (k_4)_4$ ; $f_2$ is the functional response of creditors to prey with $[f_2]_4(0) = 0, (f_2^4)_4(Q) > 0, \forall Q$ , and $\lim_{Q\to\infty}(f_2)_4(Q_4) = [k_2]$ ; $\beta$ is the age dependent birth rate with $\beta(a)$ $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 4, h^4(Q)$ $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ .

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_4 + (\rho_a)_4 = -\mu_4(a_4) + (f_4)_4P_4) + \big(f_2)_4(Q_4)\big)(\rho_4)_4 \quad (4)$$

$$\rho_4(t_4, 0) = h_4(Q_4)\int_0^\infty \beta_4(a_4)\rho_4(t, a_4)da, \quad (5)$$

$$\frac{dQ_4}{dt}$$

$$= -b_4Q_4 + g_4(P_4)Q_4 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_4$ and $Q_4$ of (4-6) satisfy

$$0 \leq \rho_4(t, a_4) \leq \rho_4^*(t, a_4),$$

$$0 \leq Q_4(t) \leq Q_4^*(t),$$

Where $\rho_4^*(t, a_4)$ and $Q_4^*(t)$, are solutions of

$$(\rho_t^*)_4 + (\rho_a^*)_4 = -(\mu_4(a_4) + (f_4)_4(P_4^*) + (f_2)_4 Q_4^*))\, \rho_4^*, \quad (7)$$

$$\rho_4^*(t, 0) = \int_0^\infty \beta_4(a_4)\, \rho_4^*(t, a_4) da_4 \quad (8)$$

$$\frac{dQ_4^*}{dt} = -b_4 Q_4^* + g_4(P_4^*)Q_4^*, \quad (9)$$

With $P_4^*(t) = \int_0^\infty \rho_4^*(t, a_4) da_4$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_4^*(t)}{dt} = P_4^*(t)\big(r_4 - (f_4)(P_4^*) - (f_2)_4(Q^*)\big), \quad (10)$$

$$\frac{dQ_4^*(t)}{dt} = Q_4^*(t)(-b_4 + g_4(P_4^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_4(a) e^{-\int_0^a \mu_4(\tau)dt} e^{-ra}\, da_4 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_4^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^{\infty} \beta_4(a_4) e^{-\int_0^a \mu_4(t)dt}\, da_4 > 1, \qquad (H1)$$

Which guarantees $r_4 > 0$.

*Remark.* If $\beta_4(a) = [\beta_0]_4 a_4 e^{-\alpha_4 a_4}, \mu_4(a_4) \equiv \mu_4$ then the solution of (12) is

$$r_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_4]_4(P_4) = 0, [f_2](Q_4) = \lambda_4 Q_4\ h_4(Q_4) = 1, \qquad and\ \ g_4(P_4) = c_4 P_4$$

$$[\rho_4]_4 + [\rho_a]_4 = -(\mu_4 + \lambda_4 Q_4)\rho_4$$

$$\rho_4(t,0) = \int_0^{\infty} [\beta_0]_4 a_4 e^{-\alpha_4 a_4} \rho_4(t, a_4)\, da_4$$

$$\frac{dQ_4}{dt} = -b_4 Q_4 + c_4 P_4 Q_4$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_4}{dt} = \omega_4 P_4 - \lambda_4 P_4 Q_4$$

$$\frac{dQ_4}{dt} = b_4 Q_4 - c_4 P_4 Q_4$$

Where $\omega_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$

## 2. THE EQUILIBRIUM $(\rho_0)_4(a_4), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_4(a_4), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_4 + (\rho_a)_4 = -\big(\mu_4(a_4) + (f_4)_4(P_4)\big)\rho,$$

$$\rho_4(t,0) = \int_0^\infty \beta_4(a_4)\rho_4(t,a_4)da_4 \qquad (13)$$

The solution of (13) has the form of

$$\rho_4(t,a_4) = Ce_4^{r(t-a_4)} - \int_0^a (\tau)d\tau - \int_0^4 (f_4)_4\big(P_4(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_4(t) = C_4 e^{r_4 t - \int_0^t (f_4)_4(P_4(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_4^{-4}(r)$ if $k_4 > r$. This is equivalent to

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{(k_4)_4 a_4} da_4 < 1.$$

In fact, if we define

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau)d\tau} e^{-s_4 a_4}\, da_4$$

$R(s)$ is a decreasing function. $R_4(r_4) = 1$ And $R_4(k_4)_4 < 1$ imply $r_4 < (k_4)_4$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_4(a_4), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{k_4 a_4} da_4 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_4 \equiv \int_0^\infty (\rho_0)_{44}(a_4)da_4$ *satisfies*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{-(f_4)_4(P_0)a_4} da = 1 \qquad (11)$$

Next, we study the stability of $(\rho_0)_4(a_4), 0)$

Let $u_4(t, a_4) = \rho(t, a_4) - \rho_0(a_4)$ Then

$$N_4(t) \equiv \int_0^\infty u_4(t, a_4)da_4 = \int_0^\infty \rho_4(t, a_4)da_4 - \int_0^\infty (\rho_0)_4(a_4)da = P_4(t) - P_0)_4$$

And

$$(u_4)_4 + (u_a)_4 = (\rho_t)_4 + (\rho_a)_4 - (\rho_0^4)_4(a_4)$$

$$= -\mu_4(a_4) + (f_4)_4(P_4) + (f_2)_4(Q_4))\,(u_4 + (\rho_0)_4) + \mu_4(a_4)(\rho_0)_4 + (f_4)_4(P_0)_4(\rho_0)_4$$

$$= -\mu_4(a_4)u_4 - (f_4)_4(P_4) + (f_2)_4(Q_4))\,u_4 - (\rho_0)_4\,(f_4)_4(P_4) - (f_4)_4(P_0)_4$$

$$+ (f_2)_4(Q_4)), \qquad (15)$$

Where $(\rho_0^4)_4(a_4)$ denotes $d(\rho_0)_4(a_4)/da_4$

Expansion $(f_4)_4(P_4)$ and $(f_2)_4(Q_4)$ at $(P_0)_4$ and 0 yield respectively

$$(f_4)_4(P_4) = (f_4)_4((P_0)_4 + N_4)$$

$$= (f_4)_4(P_0)_4 + (f_4^4)_4(P_0)_4 N_4 + o_4(N_4), \quad (16)$$

$$(f_2)_4(Q_4)$$

$$= (f_2^4)_4(0)Q + o_4(Q_4) \quad (17)$$

Write $f_4(P_0) = f_{40}, f_4'(P_0) = f_{40}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_4)_4 + (u_a)_4 \approx -(\mu_4(a_4) + (f_{40})_4)u - (f_{40}^4)_4(\rho_0)_4(N_4) - (f_{20}^4)_4(\rho_0)_4(Q_4)$$

$$u_4(t,0) = \rho_4(t,0)(\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a) + (\rho_0)_4(a_4)da_4 - (\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a_4)da_4 + (h_4(Q_4) - 4)(\rho_0)_4(0) \quad (18)$$

Expanding

$$h_4(Q_4) = 1 + (h^4)_4(0)Q_4 + 0(Q_4) \equiv 1 + h_0'Q_4 + 0Q_4$$

And substituting it into (48), we arrive at

$$u_4(t,0) \approx \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + (h_{a_4}')(\rho_0)_4 Q_4$$

The linearization is the following :

$$[u_4] + [u_4] = -(\mu_4(a_4) + [t_{40}]_4 u_4 - [\alpha_4]_4(a_4)N_4 - [\alpha_2]_4(a_4)Q_4$$

$$u_4(t,0) = \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + [\xi_4]Q_4 \quad (19)$$

$$\frac{dQ_4}{dt} = -b_4 Q_4$$

where

$$[\alpha_4]_4(a_4) = [t_4^4 a]_4[p, o]_4(a_4) d2(a_4) = [f_{20}^4]_4(\rho\, o\,)_4(a_4)\xi = [\mathrm{h\ o}^4]_4[\rho \mathrm{o}]_4(\mathrm{o}) \text{ and } \mathrm{b}_4 = \mathrm{b}_4 - g_4(P_4)$$

If $g_4(P_0) > b$, clearly, $(\rho_0)_4(a_4), 0)$ is unstable and the predator population is persistent. Assume $g_4(P_0)_4 < b$.

Then, $Q_4(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_4 + u_a = -(\mu_4(a_4) + (f_{40})_4)u_4 - (\alpha_4)_4(a_4)N_4$$

$$= -(\mu_4(a_4) + f_{40})u_4 - \alpha_4(a_4)\int_0^\infty u_4(t, a_4)da_4$$

$$u_4(t, 0) = \int_0^\infty \beta_4\,(a_4)u_4(t, a_4)da_4$$

Using the method developed by Cushing [2, 3], we suppose

$$u_4(t, a_4) = \tilde{\mathrm{u}}_4(a_4)^{e^{\gamma(t-a_4)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_4}\, da_4$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_4)$ gives,

$$\tilde{\mathrm{u}}_4(a_4) = e^{-\int_0^a \mu_4(\tau_4)d\tau - (f_{40})_4 a}\left(\tilde{\mathrm{u}}_4(0) - w\int_0^a e^{\int_0^\tau \mu_4(s_4)ds + (f_{40})}\tau_4\,(\alpha_4)_4(\tau_4)e^{\gamma_4\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1-\int_0^{\infty}\beta_4(a_4)e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\right)\tilde{u}_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a(\alpha_4)_4(\tau_4)e^{-\gamma_4\tau_4}e^{(f_{40})_4\tau_4+\int_0^\tau \mu_4(s_4)ds}d\tau_4\,da_4$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}da\right)(\tilde{u}_4(0))_4$$

$$+\left(-4-\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a e^{(f_{40})_4\tau_4+\int_0^\tau \mu_4(s_4)ds+y\tau}\alpha_4(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(t_4)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_4\,(a_4)Y_4(\gamma_4,a_4)da_4\right)u_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)\int_0^a(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4=0,$$

(24)

$$\left(\int_0^{\infty}Y_4(\gamma_4,a_4)da_4\right)\tilde{u}_4(0)$$

$$+\left(-1-\int_0^{\infty} Y_4(\gamma_4,a_4)\int_0^{a}(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4\right)w=0.$$

Equqtion (20) has a solution of the form of $u_4(t,a_4)=u_4(a_4)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_4(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_4(\gamma_4,a_4)=\int_0^a(\rho_0)_4(\tau_4)Y^{-4}(\gamma_4,a_4)d\tau_4$. We obtain the characteristic equation

$$f_{40}^{l}{}_4\left(\int_0^{\infty} Y_4(\gamma_4,a_4)da_4\right)\left(\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$+\left(1-\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)da\right)\left(1+(f_{40}^{l})_4\int_0^{\infty}Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_4(a_4)=(\rho_0)_4(0)e^{-\int_0^a\mu_4(\tau_4)d\tau_4-(f_{40})_4a_4}$$

$$I_4(\gamma_4,a_4)=\int_0^a(\rho_0)_4(\tau_4)e^{\gamma_4\tau_4+(f_{40})_4\tau+\int_0^{\tau}\mu_4(s_4)ds_4}d\tau$$

$$=\frac{(\rho_0)_4(0)}{\gamma_4}(e^{\gamma_4a_4}-1)$$

Let

$$E_4(a_4)=e^{-(f_{40})_4a_4-\int_0^a\mu_4(\tau_4)d\tau_4}$$

Then $[Y_4]_4[\gamma_4]_4, a_4) = E_4(a_4)e^{-\gamma_4 a_4}$. by a simple calculation,

$$\int_0^\infty \beta_4(a_4)Y_4(\gamma_4, a_4)I(\gamma_4, a_4)da = \frac{[\rho_0]_4(0)}{\gamma}\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}\ da\right);$$

$$\int_0^\infty Y_4(\gamma_4, a_4)I(\gamma_4, a_4)da = \frac{[\rho_0]_4(0)}{\gamma_4}\left(\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}\ da_4\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_4(0)}{\gamma_4}[f_{40}^4]\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}da_4\ (\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4)$$

$$+\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right).$$

$$\left(1 + \frac{(\rho_0)_4(0)}{\gamma_4}(f_{40}^{'})_4\left(\int_0^\infty E_4(\,a_4)da_4 - \int_0^\infty E_4(\,a_4)e^{-\gamma_4 a_4}da_4\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right)(1 + \frac{(f_{40}^{'})_4}{\gamma_4}(\rho_0)_4(0)\int_0^\infty E_4(\,a_4)da_4) = 0.$$

We need ũ(a) $\not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1 - \int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)da_4 = 1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4 \neq 0,$$

Since $\int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)\ \int_0^a \alpha_4\,(\tau_4)Y^{-4}(\gamma_4, \tau_4)d\tau_4 da_4 > 0$. Therefore, the characteristic solution is

$$\gamma_4 = -(f_{40}^{'})_4(\rho_0)_4 \int_0^{\infty} E_4(a_4)da_4.$$

Since $f_4'(P_0)_4 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_4), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_4), 0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_4 + (\mu_4(a_4) + [f_4]_4(P_4^*) + [f_2]_4(Q_4^*))P_4^* = 0, \quad (26)$$

$$p_4^*(0) = h_4(Q_4^*)\int_0^{\infty} \beta_4(a_4)p_4^*(a_J)p_4^*(a_4)da_4 \quad (27)$$

$$g_4(P_4^*) - b_4 = 0, \quad (28)$$

Where $P_4^* = \int_0^{\infty} p_4^*(a_4)\, da_4$

Equation (28) yields $P^* = g^{-4}(b)$,

From (26),

$p_4^*(a) = p_4^*(0)e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}$

Substituting into (27), we need to solve

$$\left(1 - h_4(Q_4^*)\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}da_4\right)p_4^*(0) = 0$$

For $Q_4^*$

Fix $p_4^*(0)$ and $P_4^*$ Let

$$H_4(z_4) = 1 - h_4(z_4)\int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}da_{,4}$$

And assume that

$$[k_3]_4\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(g_4^{-4}(b_4))a_4 - [k_2]_4 a_4}\, da_4 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_4(z_4)] > 0.$$

Clearly,

$$\frac{dH_4}{dz_4} = h_4(z)([f_2']_4(z_4))\int_0^\infty a_4\beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - f_4(P_4^*)a_4 - f_2(z_4)a_4}da_{\square}$$

$-h'(z_4)\int_0^\infty \beta_4(a_4)\, e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}\quad da_4$

the assumptions on $\left[f_2^4\right]_4 and\ [h^4]^4 dH_4/dz > 0$. Hence, if

$$H_4(0) = 1 - \int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4} da_4 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_4(0) < 0$ follows if $g_4^-(b_4) < [f_4^{-4}]_4(r_4)$, where $r_4$ is the solution of (12).

Indeed, since

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 e^{-y\frac{a}{4}}} da_4$$

Is a decreasing function of $s_4$ when $P_4^* = [g^{-4}]_4(b_4) < [f_4^{-4}]_4(r_4)$, that is $[f_4]_4(P_4^*) < r$, $R\left(f_4(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_4(g^{-4}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4\left(g_4^{-4}(b)\right)a_4} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_4^*(a_4), Q_4^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_4^*(a_4), Q_4^*)$ is more complicated. We again consider perturbations $m_4(t_4 a_4) = p_4(t_4, a_4) - p_4^*(a_4)$ and $W_4(t) = Q_4(t) - Q^*$.Then

$$M_4(t) \equiv \int_0^\infty m_4(t, a_4) da_4 = P_4(t) - P_4^*,$$

$$[m_4]_4 + [m_a]_4 = [p_4]_4 + [p_a]_4 - [p^*]_4(a_4) \approx -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4$$

$$-[f_{4*}^4]_4 p_4^*(a_4)M - [f_{2*}^4]_4 p_4^*(\mathrm{a})W_4,$$

$$m_4(t,0)p_4(t,0) - p_4^*(0) \approx h_* \int_0^\infty \beta_4(a_4)m_4(t,a_4)\,da + \frac{[h_*^4]_4}{[h_*]_4}\, p_4^*(0)W_4$$

$$\frac{dW_4}{dt} = \frac{dQ_4}{dt} \approx [g_*^4]_4 Q_4^* M_{,4}$$

Where $[f_{4*}]_4$ denoting $[f_4(P^*)]_4$, $[f_{4*}]_4$ denoting $[f_2(P^*)]_4$, $[f_{4*}^{'}]_4$ denoting $[f^{'}{}_4(P^*)]_4 [f_{2*}^{'}]_4$ denoting $[f_2]_4$ $(Q_4^*,$ $h_{*}$ denoting $h_*^{'}$ denoting $h_4^{'}$ $(Q_4^*)$, and $[g_*^{'}]_4$

Denoting $[g^{'}]_4(P_4^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_4]_4 + [m_a]_4 = -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4 - [f^{'}{}_{4*}]_4 [p_4^*](a_4)M_4 - [f_{2*}^{'}]_4 p_4^*(a_4)W_4$$

$$m_4(t,0) = [h_*]_4 \int_0^\infty \beta_4(a_4)m_4(t,a_4)da_4 + \frac{[h_*^4]_4}{[h^*]_4}[p^*]_4(0)\, W_{,4} \qquad (30)$$

$$\frac{dW_4}{dt} = [g_*^4]_4 [Q_4^*]M_4$$

This is equation is important in the following cases: problem of using observed correlations to infer causal relations is relevant to (e) a wide variety of scientific disciplines. Yet given correlations between just two classical variables, (eb) it is impossible to determine whether they arose from (e) a causal influence of one on the other or a common cause influencing both, unless one can implement (eb) a randomized intervention. **Katja Ried, Megan Agnew,**

**Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. Resch** here consider the problem of causal inference for (e) quantum variables. Authors introduce causal tomography, which unifies and generalizes (e&eb) conventional quantum tomography schemes to provide (eb) a complete solution to the causal inference problem using (e) a quantum analogue of a randomized trial. Authors furthermore show that, in contrast to the classical case, observed quantum correlations alone can sometimes provide (eb) a solution. They implement a quantum-optical experiment that allows (eb) to control (e&eb) the causal relation between two optical modes, and two measurement schemes -- one with (e&eb) and one without (e) randomization -- that extract this relation from (e) the observed correlations. Results show that entanglement and (e&eb) coherence, known to be central to (e) quantum information processing, also provide (eb) a quantum advantage for (e) causal inference. Subjects: Quantum Physics (quant-ph); Learning (cs.LG); General Relativity and Quantum Cosmology (gr-qc); Machine Learning (stat.ML) Journal reference: Nat Phys 11, 414-420 (2015) DOI: 10.1038/nphys3266 Cite as: arXiv: 1406.5036 [quant-ph] (or arXiv: 1406.5036v1 [quant-ph] for this version) **Inferring causal structure: a quantum advantage Katja Ried, Megan Agnew, Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. ReschSTOP D. T. Son, A. O. Starinets** review recent progress in applying the AdS/CFT correspondence to (e&eb) finite-temperature field theory. In particular, they show how the hydrodynamic behavior of field theory is reflected in (e&eb) the low-momentum limit of (e) correlation functions computed through (e&eb) a real-time AdS/CFT prescription, which they formulate. Authors also show how the hydrodynamic modes in field theory correspond to (e&eb) the low-lying quasinormal modes of (e) the AdS black p-brane metric. They provide a proof of the universality of (e) the viscosity/entropy ratio within (eb) a class of theories with (e&eb) gravity duals and formulate (eb) a viscosity bound conjecture. Possible implications for real systems are mentioned. Subjects: High Energy Physics - Theory (hep-th) Journal reference: Ann.Rev.Nucl.Part.Sci.57:95-118,2007 DOI: 10.1146/annurev.nucl.57.090506.123120Report number: INT PUB 07-02Cite as: arXiv:0704.0240 [hep-th]

(or arXiv:0704.0240v2 [hep-th] for this version) An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society **Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP** Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv:

1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version **Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP Viscosity, Black Holes, and Quantum Field Theory D. T. Son, A. O. Starinets STOP Authors** study information retrieval from evaporating black holes, assuming that the internal dynamics of a black hole is unitary and rapidly mixing, and assuming that the retriever has unlimited control over the emitted Hawking radiation. If the evaporation of the black hole has already proceeded past the ``half-way" point, where (e) half of the initial entropy has been radiated away, then (eb) additional quantum information deposited in (eb) the black hole is revealed in (eb) the Hawking radiation very rapidly. Information deposited prior to the half-way point remains concealed until (e) the half-way point, and then emerges (eb) quickly. These conclusions hold because (e) typical local quantum circuits are efficient encoders for (e) quantum error-correcting codes that nearly achieve (eb) the capacity of the quantum erasure channel. Estimate of a black hole's information retention time, based on (e) speculative dynamical assumptions, is just barely compatible with (e&eb) the black hole complementarity hypothesis. **Black holes as mirrors: quantum information in random subsystems Patrick Hayden1 and John Preskill2 Published 26 September 2007 Journal of High Energy Physics, Volume 2007, JHEP09 (2007) STOP** By analyzing the key properties of black holes from the point of (e) view of quantum information, authors derive a model-independent picture of (e) black hole quantum computing. It has been noticed that this picture exhibits (eb) striking similarities with (e&eb) quantum critical condensates, allowing (eb) the use of a common language to describe (eb) quantum computing in both systems. Authors analyze such quantum computing by allowing coupling to (e&eb) external modes, under the condition that the external influence must be (=) soft-enough in order not to (e) offset the basic properties of the system. Authors derive model-independent bounds on some crucial time-scales, such as (=) the times of gate operation, decoherence, maximal entanglement and total scrambling. They show that for black hole type quantum computers all these time-scales are of (e) the order of the

black hole half-life time. Furthermore, authors construct explicitly a set of Hamiltonians that generates (eb) a universal set of quantum gates for (e) the black hole type computer. They find that the gates work at maximal energy efficiency. Furthermore, authors establish a fundamental bound on the complexity of quantum circuits encoded on (eb&e) these systems, and characterize the unitary operations that are (=) implementable. It becomes apparent that the computational power is very limited due to (e) the fact that the black hole life-time is of the same order of (e) the gate operation time. As a consequence, it is impossible to retrieve its information, within (eb) the life-time of a black hole, by externally coupling to (e&eb) the black hole qubits. However, authors show that, in principle, coupling to (e&eb) some of the internal degrees of freedom allows (eb) acquiring knowledge about the micro-state. Still, due to the trivial complexity of operations that can be performed, (eb) there is no time advantage over the collection of Hawking radiation and subsequent decoding. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Report number: LMU-ASC 03/16, MPP-2016-1 Cite as: arXiv:1605.01407 [hep-th] (or arXiv:1605.01407v1 [hep-th] for this version) **Universality of Black Hole Quantum Computing Gia Dvali, Cesar Gomez, Dieter Lust, Yasser Omar, and Benedikt Richter (Submitted on 4 May 2016) STOP** In this addendum to [arXiv: 1402.5674] two points are discussed. In the first additional evidence is provided (eb) for a dual connection between the geometric length of an Einstein-Rosen bridge and (e&eb) the computational complexity of the quantum state of the dual CFT's. The relation between growth of complexity and (e&eb) Page's “Extreme Cosmic Censorship” principle is also remarked on. The second point involves a gedanken experiment in which Alice measures (eb) a complete set of commuting observables at her end of (e) an Einstein-Rosen bridge is discussed. An apparent paradox is resolved by appealing to (e) the properties of GHZ tripartite entanglement. **Addendum to computational complexity and black hole horizons Leonard Susskind* Version of Record online: 18 JAN 2016 DOI: 10.1002/prop.201500093 STOP** Analyzing some well established facts, authors give a model-independent

parameterization of black hole quantum computing in terms of (e&*eb) a set of macro and micro quantities and their relations. These include the relations between the extraordinarily-small energy gap of (e) black hole qubits and important time-scales of (e) information-processing, such as, scrambling time and Page's time. Authors then show, confirming and extending previous results, that other systems of nature with (e&eb) identical quantum informatics features are (=) attractive Bose-Einstein systems at (eb) the critical point of quantum phase transition. Here authors establish a complete isomorphy between (e&eb) the quantum computational properties of these two systems. In particular, they show that the quantum hair of a critical condensate is (=) strikingly similar to (=) the quantum hair of a black hole. Irrespectively whether one takes the similarity between the two systems as a remarkable coincidence or as (=) a sign of a deeper underlying connection, the following is evident. Black holes are not unique in their way of (e) quantum information processing and they can manufacture black hole based quantum computers in labs by (e) taking advantage of quantum criticality. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Atomic Physics (physics. Atom-ph); Quantum Physics (quant-ph) Cite as: arXiv: 1601.01329 [hep-th] (or arXiv: 1601.01329v1 [hep-th] for this version) **Black Hole Based Quantum Computing in Labs and in the Sky Gia Dvali, Mischa Panchenko (Submitted on 6 Jan 2016) STOP Gia Dvali, Alexander Gußmann** show that the existence of black holes with (e&eb) classical skyrmion hair invalidates (e) standard proofs that global charges, such as (=) the baryon number, cannot be conserved by (e) a black hole. By carefully analyzing the standard arguments based on (e) a Gedankenexperiment in which a black hole is seemingly-unable to (e) return the baryon number that it swallowed, authors identify inconsistencies in this reasoning, which does not (e) take into the account neither the existence of skyrmion black holes nor the baryon/skyrmion correspondence. Authors then perform a refined Gedankenexperiment by incorporating (e) the new knowledge and show that (eb) no contradiction with conservation of baryon number takes place at (eb) any stage of black hole evolution. Analysis also indicates (eb) no conflict

between semi-classical black holes and (e&eb) the existence of baryonic gauge interaction arbitrarily-weaker than gravity. Next, authors study classical cross sections of a minimally-coupled massless probe scalar field scattered by (e) _ a skyrmion black hole. Authors investigate how the skyrmion hair manifests itself by (e) comparing this cross section with (e&eb) the analogous cross section caused by (e) a Schwarzschild black hole which has (e) the same ADM mass as (=) the skyrmion black hole. Here authors find an order-one difference in the positions of (e) the characteristic peaks in (eb) the cross sections. The peaks are shifted to smaller scattering angles when (e) the skyrmion hair is present. This comes from the fact that the skyrmion hair changes (e&eb) the near horizon geometry of the black hole when compared to (e&eb) a Schwarzschild black hole with same ADM mass. Authors keep the study of this second aspect general so that (e) the qualitative results obtained can also be applied to(e&eb) black holes with classical hair of different kind. Subjects: High Energy Physics - Theory (hep-th); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU-ASC 18/16 Cite as: arXiv:1605.00543 [hep-th] (or arXiv:1605.00543v2 [hep-th] for this version) A preliminary discussion is given of (e) the prospects that gravitational-wave observations of (e) binary inspiral of black holes could reveal or constrain (e&eb) quantum modifications to (e) black hole dynamics, such as are required to (e) preserve postulates of quantum mechanics. Different proposals for such modifications are characterized by (e) different scales, and the size of these scales relative to (e&eb) those probed by observation of (e) inspiral signals is important in determining (eb) the feasibility of finding experimental signatures. Certain scenarios with strong quantum modifications in (eb) a region extending well outside the horizon are expected to modify (e&eb) classical evolution, and distort (e&eb) the near-peak gravitational wave signal, suggesting (eb) a search for departures from (e) waveforms predicted by general relativity. The near agreement of the GW150914 signals with (e&eb) such waveforms is discussed, and indicates (eb) constraints on some such scenarios. Important strategies for more precise future tests are 1) to develop more precise

predictions from (e) scenarios proposing quantum modifications, and 2) searching for observed deviations from (e) numerical relativity predictions via (e&eb) analysis of gravity wave data, particularly focussing on (e&eb) the signal region corresponding to plunge and merger. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1602.03622 [gr-qc] (or arXiv: 1602.03622v2 [gr-qc] for this version) **Gravitational wave tests of quantum modifications to black hole structure -- with post-GW150914 update Steven B. Giddings STOP** Direct LHC bounds on (e&eb) colored SUSY particles now corner naturalness more than the measured value of the Higgs mass does. Bounds on the gluino are of (e) particular importance, since it radiatively "sucks" up (e) the stop and Higgs-up soft masses. As a result, even models that easily accommodate a 125 GeV Higgs are almost as tuned as (=) the simplest version of SUSY, the MSSM: at best at the percent level. In this paper, authors further examine how current LHC results constrain (e) naturalness in three classes of models that may relax LHC bounds on (e&eb) sparticles: split families (e&eb), baryonic RPV (e&eb), and Dirac gauginos. In models of split families and bRPV, the bounds on the gluino are (=) only slightly reduced, resulting in (eb) a few percent tuning. In particular, having a natural spectrum in bRPV models typically implies (eb) that tops, Ws, and Zs are easily produced (eb) in the cascade decays of (e) squarks and gluinos. The resulting leptons and missing energy push (eb) the gluino mass limit above 1 TeV. Even when the gluino has a Dirac mass and does not contribute to the stop mass at one loop, tuning reappears in calculable models because there is no symmetry imposing the supersoft limit. Authors conclude that, even if sparticles are found at LHC-14, naturalness will not emerge triumphant. Subjects: High Energy Physics - Phenomenology (hep-ph)DOI:10.1007/JHEP03(2014)022 Cite as:rXiv:1309.3568 [hep-ph] (or arXiv:1309.3568v2 [hep-ph] for this version) **The Last Vestiges of Naturalness Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang, Ken Van Tilburg, Giovanni Villadoro STOP** The lack of evidence for new physics beyond (e) the standard model at the LHC points to (e) a paucity of new particles near the weak scale. This suggests that the weak scale is tuned and that supersymmetry, if present at

all, is realized (eb) at higher energies. The measured Higgs mass constrains (e) the scalar sparticles to be below 10^5 TeV, while (e) gauge coupling unification favors (eb) Higgsinos below 100 TeV. Nevertheless, in many models gaugino masses are (=) suppressed and remain within reach of the LHC. Tuning the weak scale and (e&eb) the renormalization group evolution of (e) the scalar masses constrain (e) Split model building. Due to the small gaugino masses, either the squarks or the up-higgs often run tachyonic; in the latter case, successful electroweak breaking requires heavy higgsinos near the scalar sparticles. We discuss the consequences of tuning the weak scale and the phenomenology of several models of Split supersymmetry including anomaly mediation, U(1)_(B-L) mediation, and Split gauge mediation. Subjects: High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Experiment (hep-ex) DOI: 10.1007/JHEP02 (2013)126 Cite as: arXiv: 1210.0555 [hep-ph] (or arXiv: 1210.0555v2 [hep-ph] for this version) **Mini-Split Asimina Arvanitaki, Nathaniel Craig, Savas Dimopoulos, Giovanni Villadoro STOP** Light states associated with (e&eb) the hierarchy problem affect (e&eb) the Higgs LHC production and decays. Authors illustrate this within the MSSM and two simple extensions applying (e&eb) the latest bounds from (e) LHC Higgs searches. Large deviations in the Higgs properties are expected in (eb) a natural SUSY spectrum. The discovery of a non-Standard-Model Higgs may signal (eb) the presence of light stops accessible at (eb) the LHC. Conversely, the more the Higgs is Standard-Model-like, (eb) the more tuned the theory becomes. Taking the ratio of different Higgs decay channels at the LHC cancels (er&eb) the leading (eb) QCD uncertainties and potentially improves (eb) the accuracy in Higgs coupling measurements to (e) the percent level. This may lead to (eb) the possibility of doing precision Higgs physics at (eb) the LHC. Finally, authors entertain the possibility that the ATLAS excess around (e&eb) 125 GeV persists with (e&eb) a Higgs production cross-section that is enhanced compared to (e&eb) the SM. This increase can only be accommodated in extensions of (e) the MSSM and it may suggest (eb) that stops lie below 400 GeV, likely within (eb) reach of next year's LHC run. Subjects: High Energy Physics - Phenomenology (hep-ph) DOI: 10.1007/JHEP02 (2012)144 Cite

as: arXiv: 1112.4835 [hep-ph] (or arXiv: 1112.4835v2 [hep-ph] for this version) **A Non Standard Model Higgs at the LHC as a Sign of Naturalness Asimina Arvanitaki, Giovanni Villadoro STOP** It has recently been suggested that the presence of a plenitude of light axions, an Axiverse, is (=) evidence for the extra dimensions of string theory. Authors discuss the observational consequences of (e) these axions on astrophysical black holes through (e&eb) the Penrose superradiance process. When an axion Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e&eb) the black hole "nucleus" forming (eb) a gravitational atom in the sky. The occupation number of superradiant atomic levels fed by (e) the energy and angular momentum of the black hole grows exponentially. The black hole spins down and an axion Bose-Einstein condensate cloud forms (eb) around it. When the attractive axion self-interactions become (=) stronger than the gravitational binding energy, (eb) the axion cloud collapses, a phenomenon known in condensed matter physics as "Bosenova". The existence of axions is first diagnosed by (e) gaps in the mass vs (e&eb) spin plot of astrophysical black holes. For young black holes the allowed values of (e) spin are quantized, giving rise to (eb) "Regge trajectories" inside the gap region. The axion cloud can also be observed directly either through (e&eb) precision mapping of the near horizon geometry or through (e&eb) gravitational waves coming from (e) the Bosenova explosion, as well as axion transitions and (e&eb) annihilations in (eb) the gravitational atom. Estimates suggest (eb) that these signals are detectable in (eb) upcoming experiments, such as Advanced LIGO, AGIS, and LISA. Current black hole spin measurements imply (eb) an upper bound on (e&eb) the QCD axion decay constant of (e) 2 x 10^17 GeV, while (e) Advanced LIGO can detect (e&eb) signals from (e) a QCD axion cloud with (e&eb) a decay constant as low as (=) the GUT scale. Authors finally discuss the possibility of observing the gamma-rays associated with the Bosenova explosion and, perhaps, the radio waves from axion-to-photon conversion for the QCD axion. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Journal

reference:Phys.Rev.D83:044026, 2011 DOI: 10.1103/PhysRevD.83.044026 Cite as: arXiv: 1004.3558 [hep-th] (or arXiv: 1004.3558v2 [hep-th] for this version) **Exploring the String Axiverse with Precision Black Hole Physics Asimina Arvanitaki, Sergei Dubovsky (Submitted on 20 Apr 2010 (v1), last revised 16 Aug 2011 (this version, v2)) STOP Skyrmion Black Hole Hair: Conservation of Baryon Number by Black Holes and Observable Manifestations Gia Dvali, Alexander Gußmann STOP Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust** discuss BMS supertranslations both at (eb) null-infinity and on the horizon for the case of the Schwarzschild black hole. Authors show that both kinds of supertranslations lead to (eb) infinitely many gapless physical excitations. On this basis authors construct a quotient algebra using (e) suited superpositions of (e) both kinds of transformations which cannot be compensated by (e) an ordinary BMS-supertranslation and therefore are (=) intrinsically due to (e) the presence of an event horizon. Authors show that these quotient transformations are (=) physical and generate (eb) gapless excitations on the horizon that can account for (e) the gravitational hair as well as for (e) the black hole entropy. They identify the physics of these modes as associated with (e&eb) Bogolioubov-Goldstone modes due to quantum criticality. Classically the number of these gapless modes is (=) infinite. However, authors show that due to (e) quantum criticality the actual amount of information-carriers becomes (=) finite and consistent with Bekenstein entropy. Although they only consider the case of Schwarzschild geometry, the arguments are extendable to(e&eb) arbitrary space-times containing event horizons. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU--ASC 05/16, MPP--2016—3 Cite as: arXiv: 1601.03725 [hep-th] (or arXiv: 1601.03725v1 [hep-th] for this version) **Gravitational Black Hole Hair from Event Horizon Supertranslations Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust STOP Gia Dvali, Cesar Gomez, Nico Wintergerst** suggest that holography can be formulated in terms of (e&eb) the information capacity of (e) the St\"uckelberg degrees of freedom that maintain

gauge invariance of (e) the theory in the presence of (e) an information boundary. These St\"uckelbergs act as (e&eb) qubits that account for a certain fraction of (e) quantum information. Their information capacity is measured by (e) the ratio of the inverse St\"uckelberg energy gap to (e) the size of the system. Systems with (e&eb) the smallest gap are (=) maximally holographic. For massless gauge systems this information measure is (=) universally equal to (=) the inverse coupling evaluated at the systems' length scale. In this language it becomes very transparent why the St\"uckelberg information capacity of black holes saturates (e full) the Bekenstein bound and accounts for (e) the entire information of the system. The physical reason is that the strength of quantum interaction is bounded from below by the gravitational coupling, which scales as (=) area. Observing the striking similarity between the scalings of the energy gap of the boundary St\"uckelberg modes and (e&eb): (AC=DC=0) the Bogoliubov modes of critical many-body systems, authors establish a connection between holography and (e&eb) quantum criticality through (e&eb) the correspondence between these modes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1511.03525 [hep-th] (or arXiv: 1511.03525v1 [hep-th] for this version) **Stückelberg Formulation of Holography Gia Dvali, Cesar Gomez, Nico Wintergerst STOP B**lack hole entropy, denoted by N, in (semi)classical limit is (=) infinite. This scaling reveals (eb) a very important information about the qubit degrees of (e) freedom that carry black hole entropy. Namely, the multiplicity of qubits scales as N, whereas (e) their energy gap and their coupling as 1/N. Such a behavior is indeed exhibited by (e) Bogoliubov-Goldstone degrees of freedom of (e) a quantum-critical state of N soft gravitons (a condensate or a coherent state) describing (eb) the black hole quantum portrait. They can be viewed as (=) the Goldstone modes of (e) a broken symmetry acting on (e&eb) the graviton condensate. In this picture Minkowski space naturally emerges as (=) a coherent state of infinite-N gravitons of (e) infinite wavelength and it carries (e&eb) an infinite entropy. In this paper authors **Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015)** ask what is the geometric meaning (if any) of (e) the classical limit of this symmetry. Authors argue that the infinite-

N limit of Bogoliubov-Goldstone modes of critical graviton condensate is described by (e) recently-discussed classical BMS super-translations broken by the black hole geometry. However, the full black hole information can only be recovered for (e) finite N, since the recovery time becomes (=) infinite in classical limit in which N is (=) infinite. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) DOI: 10.1016/j.physletb.2015.11.073 Report number: MPP-2015-215, LMU-ASC 59/15 Cite as: arXiv: 1509.02114 [hep-th] (or arXiv: 1509.02114v1 [hep-th] for this version) **Classical Limit of Black Hole Quantum N-Portrait and BMS Symmetry Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015) STOP** Recent ideas about understanding physics of black hole information-processing in terms of (e&eb) quantum criticality allow us to (eb) implement black hole mechanisms of quantum computing within (eb) critical Bose-Einstein systems. The generic feature, uncovered both by analytic and numeric studies, is (=) the emergence at the critical point of (e) gapless weakly-interacting modes, which act as (=) qubits for information-storage at a very low energy cost. These modes can be effectively described in terms of (e&eb) either Bogoliubov or Goldstone degrees of freedom. The ground-state at the critical point is (=) maximally entangled and far from being classical. We confirm this near-critical behavior by a new analytic method. **Gia Dvali, Mischa Panchenko** compute growth of entanglement and show its consistency with (e&eb) black hole type behavior. On the other hand, in the over-critical regime the system develops (eb) a Lyapunov exponent and scrambles quantum information very fast. By, manipulating (e&eb) the system parameters externally, authors can put it in and out of various regimes and in this way control (e&eb) the sequence of information storage and processing. By using gapless Bogoliubov modes as (=) control qubits they design some simple logic gates. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Cite as: arXiv: 1507.08952 [hep-th] (or arXiv: 1507.08952v1 [hep-th] for this version) **Black Hole Type Quantum**

**Computing in Critical Bose-Einstein Systems Gia Dvali, Mischa Panchenko (Submitted on 31 Jul 2015) STOP Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst** establish a fundamental connection between quantum criticality of a many-body system, such as Bose-Einstein condensates, and (e&eb) its capacity of information-storage and processing. For deriving the effective theory of modes in the vicinity of the quantum critical point they develop a new method by mapping a Bose-Einstein condensate of N-particles onto (e&eb) a sigma model with a continuous global (pseudo)symmetry that mixes (e&eb) bosons of different momenta. The Bogolyubov modes of the condensate are mapped onto (e&eb) the Goldstone modes of the sigma model, which become (=) gapless at the critical point. These gapless Goldstone modes are (=) the quantum carriers of information and entropy. Analyzing their effective theory, authors observe the information-processing properties strikingly similar to (=) the ones predicted by the black hole portrait. The energy cost per qubit of information-storage vanishes in (eb) the large-N limit and the total information-storage capacity increases with (e&eb) N either exponentially or as a power law. The longevity of information-storage also increases with (e&eb) N, whereas (e) the scrambling time in (eb) the over-critical regime is controlled by (e&eb) the Lyapunov exponent and scales logarithmically with (e&eb) N. This connection reveals (eb) that the origin of black hole information storage lies in (eb) the quantum criticality of the graviton Bose-gas, and that much simpler systems that can be manufactured in table-top experiments can exhibit (eb) very similar information-processing dynamics. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 92, 125002 (2015) DOI: 10.1103/PhysRevD.92.125002 Cite as: arXiv:1507.02948 [hep-th] (or arXiv:1507.02948v1 [hep-th] for this version) **Nambu-Goldstone Effective Theory of Information at Quantum Criticality Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst STOP Gia Dvali, Cesar Gomez** show that a positive cosmological constant is (=) incompatible with the quantum-corpuscular resolution of de Sitter metric in form of (e) a coherent state. The reason is very general and is due to (e) the

quantum self-destruction of (e) the coherent state because of (e) the scattering of constituent graviton quanta. This process creates an irreversible quantum clock, which precludes (e) eternal de Sitter. It also eliminates (e) the possibility of Boltzmann brains and Poincare recurrences. This effect is expected to be (=) part of any microscopic theory that takes into account (e) the quantum corpuscular structure of the cosmological background. This observation puts the cosmological constant problem in a very different light, promoting it, from a naturalness problem, into (e&eb) a question of quantum consistency. We are learning that quantum gravity cannot tolerate (e) exceedingly-classical sources. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1412.8077 [hep-th]

(or arXiv:1412.8077v1 [hep-th] for this version) **Quantum Exclusion of Positive Cosmological Constant? Gia Dvali, Cesar Gomez (Submitted on 27 Dec 2014) STOP G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger** establish a connection between the ultra-Planckian scattering amplitudes in field and string theory and (e&eb) unitarization by black hole formation in these scattering processes. **Using as a guideline an explicit microscopic theory in which the black hole represents (eb) a bound-state of many soft gravitons at (eb) the quantum critical point, they were able to identify and compute (e&eb) a set of perturbative amplitudes relevant for (e) black hole formation. These are the tree-level N-graviton scattering S-matrix elements in (eb) a kinematical regime (called classicalization limit) where (e) the two incoming ultra-Planckian gravitons produce (eb) a large number N of soft gravitons. Authors compute these amplitudes by using (e) the Kawai-Lewellen-Tye relations, as well as scattering equations and (e&eb) string theory techniques. Authors discover that this limit reveals (eb) the key features of the microscopic corpuscular black hole N-portrait. In particular, the perturbative suppression factor of a N-graviton final state, derived from (e) the amplitude, matches the non-perturbative black hole entropy when (e) N reaches the quantum criticality value, whereas final states with different value of N are (=) either suppressed or excluded by non-perturbative**

**corpuscular physics. Thus they identify the microscopic reason behind the black hole dominance over other final states including (e) non-black hole classical object. In the parameterization of the classicalization limit the scattering equations can be solved exactly allowing (eb) us to obtain closed expressions for the high-energy limit of the open and closed superstring tree-level scattering amplitudes for a generic number N of external legs. Authors demonstrate matching and complementarity between the string theory and (e&eb) field theory** in different large-s and large-N regimes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1016/j.nuclphysb.2015.02.004 Report number: MPP-2014-320, LMU-ASC 52/14 Cite as: arXiv:1409.7405 [hep-th] (or arXiv:1409.7405v2 [hep-th] for this version) **Black Hole Formation and Classicalization in Ultra-Planckian 2 -> N Scattering G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger STOP** Gravitational backgrounds, such as black holes, AdS, de Sitter and inflationary universes, should be viewed as (=) composite of N soft constituent gravitons. It then follows that such systems are close to (e) quantum criticality of graviton Bose-gas to Bose-liquid transition. Generic properties of the ordinary metric description, including (e) geodesic motion or particle-creation in the background metric, emerge as (=) the large-N limit of quantum scattering of (e) constituent longitudinal gravitons. **Gia Dvali, Cesar Gomez** show that this picture correctly accounts for (e) physics of large and small black holes in AdS, as well as reproduces (eb) well-known inflationary predictions for cosmological parameters. However, it anticipates new effects not captured by (e) the standard semi-classical treatment. In particular, **Gia Dvali, Cesar Gomez** predict observable corrections that are sensitive to (e&eb) the inflationary history way beyond last 60 e-foldings. They derive an absolute upper bound on the number of e-foldings, beyond (e) which neither de Sitter nor inflationary Universe can be approximated by (e) a semi-classical metric. However, they could in principle persist in a new type of (e) it quantum eternity state. Authors discuss implications of this phenomenon for (e&eb) the cosmological constant problem. Subjects: High Energy Physics - Theory (hep-th); Cosmology and

Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Quantum Physics (quant-ph) DOI: 10.1088/1475-7516/2014/01/023 Cite as: arXiv: 1312.4795 [hep-th] (or arXiv:1312.4795v1 [hep-th] for this version) **Quantum Compositeness of Gravity: Black Holes, AdS and Inflation Gia Dvali, Cesar Gomez STOP** As quantum information science approaches the goal of constructing quantum computers, understanding loss of information through (e&eb) decoherence becomes (=) increasingly important. The information about a system that can be obtained from (e) its environment can facilitate quantum control and error correction. Moreover, observers gain (eb) most of their information indirectly, by (e) monitoring (primarily photon) environments of (e) the "objects of interest." Exactly how this information is inscribed in the environment is essential for (e) the emergence of "the classical" from (e) the quantum substrate. In this paper, **Robin Blume-Kohout, Wojciech H. Zurek** examine how many-qubit (or many-spin) environments can store (e) information about a single system. The information lost to (e) the environment can be stored redundantly, or it can be encoded in (eb) entangled modes of the environment. Authors go on to show that randomly chosen states of the environment almost always encode (e) the information so that (e) an observer must capture (e) a majority of the environment to deduce (eb) the system's state. Conversely, in the states produced by (e) a typical decoherence process, information about a particular observable of the system is stored redundantly. This selective proliferation of "the fittest information" (known as Quantum Darwinism) plays a key role in (eb) choosing the preferred, effectively classical observables of (e) macroscopic systems. The developing appreciation that the environment functions not just as (=) a garbage dump, but as a communication channel, is extending understanding of the environment's role in (e&eb) the quantum-classical transition beyond (e) the traditional paradigm of decoherence. Keywords Decoherence quantum Darwinism quantum information quantum-classical transition **Foundations of Physics November 2005, Volume 35, Issue 11, pp 1857-1876 A Simple Example of "Quantum Darwinism": Redundant Information Storage in Many-Spin Environments Robin Blume-Kohout, Wojciech H. Zurek STOP Harold Ollivier, David**

**Poulin, and Wojciech H. Zurek** study the emergence of objective properties in open quantum systems. In our analysis, the environment is promoted from a passive role of a reservoir selectively destroying quantum coherence to an active role of amplifier selectively **proliferating information about the system. We show that only preferred pointer states of the system can leave a redundant and therefore easily detectable imprint on the environment. Observers who—as is almost always the case—discover the state of the system indirectly (by probing a fraction of its environment) will find out only about the corresponding pointer observable. Many observers can act in this fashion independently and without perturbing the system. They will agree about its state. In this operational sense, preferred pointer states exist objectively. Received 24 July 2003 DOI:http://dx.doi.org/10.1103/PhysRevLett.93.220401 Objective Properties from Subjective Quantum States: Environment as a Witness Harold Ollivier, David Poulin, and Wojciech H. Zurek Phys. Rev. Lett. 93, 220401 – Published 22 November 2004 STOP** Quantum Darwinism recognizes that we—the observers—acquire our information about the "systems of interest" indirectly from their imprints on the environment. **Michael Zwolak, H. T. Quan, and Wojciech H. Zurek** show that information about a system can be acquired from a mixed-state, or hazy, environment, but the storage capacity of an environment fragment is suppressed by its initial entropy. In the case of good decoherence, the mutual information between the system and the fragment is given solely by the fragment's entropy increase. For fairly mixed environments, this means a reduction by a factor 1−h, where h is the haziness of the environment, i.e., the initial entropy of an environment qubit. Thus, even such hazy environments eventually reveal the state of the system, although now the intercepted environment fragment must be larger by ~ (1−h)−1 to gain the same information about the system. DOI:http://dx.doi.org/10.1103/PhysRevLett.103.110402 **Quantum Darwinism in a Mixed Environment Michael Zwolak, H. T. Quan, and Wojciech H. Zurek Phys. Rev. Lett. 103, 110402 – Published 8 September 2009** We study quantum Darwinism—the redundant recording of information about the preferred states of a decohering system by its environment—for an object illuminated by a

blackbody. In the cases of point-source and isotropic illumination, we calculate the quantum mutual information between the object and its photon environment. We demonstrate that this realistic model exhibits fast and extensive proliferation of information about the object into the environment and results in redundancies orders of magnitude larger than the exactly soluble models considered to date. DOI:http://dx.doi.org/10.1103/PhysRevLett.105.020404 **Quantum Darwinism in an Everyday Environment: Huge Redundancy in Scattered Photons C. Jess Riedel and Wojciech H. Zurek Phys. Rev. Lett. 105, 020404 – Published 9 July 2010 STOP** Measurements transfer information about a system to the apparatus and then, further on, to observers and (often inadvertently) to the environment. I show that even imperfect copying essential in such situations restricts possible unperturbed outcomes to an orthogonal subset of all possible states of the system, thus breaking the unitary symmetry of its Hilbert space implied by the quantum superposition principle. Preferred outcome states emerge as a result. They provide a framework for "wave-packet collapse," designating terminal points of quantum jumps and defining the measured observable by specifying its eigenstates. In quantum Darwinism, they are the progenitors of multiple copies spread throughout the environment—the fittest quantum states that not only survive decoherence, but subvert the environment into carrying information about them—into becoming a witness. DOI:http://dx.doi.org/10.1103/PhysRevA.76.052110 **Quantum origin of quantum jumps: Breaking of unitary symmetry induced by information transfer in the transition from quantum to classical Wojciech Hubert Zurek Phys. Rev. A 76, 052110 – Published 16 November 2007 STOP** The Born rule, a foundational axiom used to deduce probabilities of events from **wavefunctions, is indispensable in the everyday practice of quantum physics. It is also key in the quest to reconcile the ostensibly inconsistent laws of the quantum and classical realms, as it confers physical significance to reduced density matrices, the essential tools of decoherence theory. Following Bohr's Copenhagen interpretation, textbooks postulate the Born rule outright. However, recent attempts to derive it from (e) other quantum principles have been successful, holding promise for (e)**

**simplifying and clarifying the quantum foundational bedrock. A major family of derivations is based on envariance, a recently discovered symmetry of (e) entangled quantum states. Here, authors identify and experimentally test three premises central to (e) these envariance-based derivations, thus demonstrating, (eb) in the microworld, the symmetries from (e) which the Born rule is derived. Further, they demonstrate envariance in a purely local quantum system, showing (eb) its independence from relativistic causality.Subjects: Quantum Physics (quant-ph)Cite as:arXiv:1604.01471 [quant-ph] (or arXiv:1604.01471v1 [quant-ph] for this version) Quantum probabilities from quantum entanglement: Experimentally unpacking the Born rule Jérémie Harris, Frédéric Bouchard, Enrico Santamato, Wojciech H. Zurek, Robert W. Boyd, Ebrahim Karimi**

Integrating $[m_4]_4(t, a_4)$ with respect to a from 0 to ∞, we have

$$\frac{dM_4(t)}{dt} = \int_0^\infty [mt]_4(t, a_4)da_4$$

$$= -\int_0^\infty [m_a]_4\,(t, a_4)da_4 - \int_0^\infty \mu(a_4)m(t, a_4)da_4$$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f^{'}{}_{4*}]_4 \int_0^\infty p_4^*(a_4)da_4\right)M_4 - [f_{2*}^4]_4 \int_0^\infty p_4^*(a_4)da_4 W_4$$

$= m_4(t, 0) - \int_0^\infty \mu_4(a_4)\, m_4(t, a_4)da_4$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}^{i}]_4 P_4^*\right)M_4 - [f_{2*}^{i}]_4 P_4^* W_{\cdot 4}$$

From (30), we arrive at the equations

$$\frac{dW_4(t)}{dt} = \int_0^{\infty} ([h_*]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4$$

$$-\left([f_{4*}] + [f_{2*}]_4 + \left[f_{4*}^{'}\right]_4 P_4^*\right) M_4 + \left(\frac{\left[h_*^{'}\right]_4}{[h_*]^4}\, P_4^*(0) - f_{2*}^{'} P_4^*\right) W_4,$$

(34)

$$\frac{dW_4(t)}{dt}\left[g_{*}^{'}\right]_4 Q_4^*\, M_{.4}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{r*}]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4 = ([h_*]_4\beta_4(\theta_4) - \mu_4(\theta_4))M_{,4}$$

Where $0< \theta_4 = \theta_4(t) < \infty$.

And if $a_4^2\,[a_4^2]_4<4\,[a_2]_4<1\, a_2 e$,

$\mathrm{Re}\,\lambda_4\left[\lambda_4\right]_4 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_4 = -\frac{4}{2}a_4[a2\ ]_4 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_*\,[h_*]_4\beta_4(a_4) - \mu_4(a_4)\right|$. Then$|n_4(t)|$ ⊠⊿∀t ⊠. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_4^2 <4\, a_2 e$ and $[a_4] > 2\Delta$,or $[a_4^2]$ ⊠[$a_2$]e and $(a_4 - \sqrt{(a_4^2)_4 - 4(a_2)_4 e}) > 2\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+4}\right)$ and $\beta_4(a_4) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4) + d\tau_4 da} = \frac{4}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_4 + [\mu_0]$ 0.5. Assume$[f_4](x_4)= K_4 x_4$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_4(x_4)= [0x_4]/(1 + x_4)$; the functional responses $[[f_2](x) = [0]\, x_4/(1 + x_4)$and $h_4(x_4)=(1+\omega_4 x_4)/(1 + x_4)$(See Holling [7] and Rashevsky [13])

Equation (44),now has the form of

$$\frac{m_4 + [\mu_0] + [K][P_0])}{[m_4]_4 + (\mu_0)_4 + [K][P_0])\,{}_4^3\,[K]_4[P_0]_4^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(4 + P_0) < b_4$ i e $[P_0] < b/(\delta_4 - b_4)$then the equilibrium$(P_0)$ $[P_0]_4(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_4$

$[f_{4*}] = K_4\, P^*\, [f_{2*}] = \frac{\sigma Q_4^*}{4+Q_4^*} \quad [h_*] = \frac{4+\omega_4 Q_4^*}{4+Q_4^*}$

$[f_{4*}] = K_4\, [f_{2*}^4] = \frac{\sigma_4}{(4+Q_4^*)^2}\, [g_*^4] = \frac{\sigma_4}{(4+P_4^*)^2} \quad [h_*^4] = \frac{\omega_4 - 4}{(4+Q_4^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0$, K, b, $\delta$, and $\omega$ need to satisfy

$$\frac{\omega_4}{(m + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma)^2} + \frac{2\omega_4}{([m_4]_4 + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma_4)^3} < 1$$

And

$$\frac{1}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^2}+\frac{2\omega_4}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4}+(m_4+[\mu_0]_4+\frac{K_4b_4}{\delta_4-b_4}+\sigma_4)\,Q_4^*)^3$$

$$=m_4+[\mu_0]_4+3+\frac{K_4b_4}{\delta_4-b_4}(4+Q^*)(1+\omega_4Q_4^*$$

Now

$$\left[a_4^2\right]-1\left[a_2\right]e=\left(\frac{2K_4b_4}{\delta_4-b_4}\right)^2+\frac{1K_4b_4\sigma_4\,Q_4^*}{(\delta_4-b_4)(1+Q_4^*)}+\frac{\sigma_4^2Q_4^{*2}}{(1+Q_4^*)^2}$$

$$\frac{1\,b_4\sigma_4\,(\delta_4-b)Q^*}{\delta_4(1+Q_4^*)^2}-\frac{1\,(\delta_4-b_4)^2(1-\omega_4)Q_4^*}{\delta_4(1+\omega_4Q_4^*)(1+Q_4^*)}P^*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_4\,ae_4^{mu}\;[\mu_0]_4-\frac{a_4}{a_4+4}$$

Hence if $[a_4^2>4(a_2)_4e$ and

$$[a_4] = \frac{2K_4 b_4}{\delta_4 - b_4} + \frac{\sigma_4 \, Q_4^*}{1 + Q_4^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_4]_4 = [f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}] P_4^*$$

$$[a_2]_4 = \left[f_{2*}^4\right] P_4^* - \frac{h_*^4}{h_*} P^*(0)$$

$$e = [g_*^4]_4 \, Q_4^*$$

$$\eta_4(t) = \left[h_*\right]_4 \beta_4 \left(\theta_4(t)\right) - \mu_4(\theta(t)),$$

The system(34)can be written as the following system:

$$[M(t)]\left[a_4\right] M_4 - [a_2]_4 W_4 - [n(t)] M_4$$

$$w_4(t) = eM_4$$

Write

$$A = \begin{pmatrix} -[a_4]_4 - [a_2]_4 & \\ e & 0 \end{pmatrix}$$

$$B = (t) = \begin{pmatrix} [\eta(t)]_4 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_4 \\ W_4 \end{pmatrix}$$

The solution of (32)is

$$X_4(t) = e^{A,t}\,[X]_4(0) + \int_0 e^{At}$$

**Some of the other conceptual and thematic and discursive factorial representationalities to which the model is applicable are:every one-dimensional quantum mechanical Hamiltonian H1 can have (e) a partner H2 such that H1 and H2 taken together may be viewed as (=) the components of a supersymmetric Hamiltonian. The term 'supersymmetric Hamiltonian' is taken to mean (eb) a Hamiltonian defined in terms of (e&eb)charges that obey (eb)the same algebra as that of (e) the generators of supersymmetry in field theory. The consequences of this symmetry for (er) the spectra of H1 and H2 are explored. It is shown how the supersymmetric pairing may be utilised to eliminate (e) the ground state of H1, or add a state below the ground state of H1 or maintain the spectrum of H1. It is also explicitly demonstrated that the supersymmetric pairing may be used to generate (eb) a class of anharmonic potentials with (e&eb) exactly specified spectra. The complete spectrum of an anharmonic potential so generated consists of (e) all the eigenstates of the simple harmonic oscillator and, in addition, a ground state at a specified energy E which lies arbitrarily below (eb) the E=1/2 ground state of the harmonic oscillator. Supersymmetric quantum mechanics of one-dimensional systems C V Sukumar Journal of Physics A: Mathematical and General, Volume 18, Number 15 STOP Authors examine the properties and consequences of (e&eb) pseudo-supersymmetry for (e) quantum systems admitting (e) a pseudo-Hermitian Hamiltonian. Authors explore the Witten index of (e) pseudo-**

**supersymmetry and show that every pair of diagonalizable (not necessarily Hermitian) Hamiltonians with (e&eb) discrete spectra and (e&eb) real or complex-conjugate pairs of (e) eigenvalues are (=) isospectral and have (e) identical degeneracy structure except (e) perhaps for the zero eigenvalue if and only if (e) they are pseudo-supersymmetric partners. This implies (eb) that pseudo-supersymmetry is (=) the basic framework for (e) generating non-Hermitian PT-symmetric and (e&eb) non-PT-symmetric Hamiltonians with (e&eb) a real spectrum via (e&eb) a Darboux transformation, and shows (eb) that every diagonalizable Hamiltonian H with(e&eb) a discrete spectrum and (e&eb) real or complex-conjugate pairs of (e) eigenvalues may be factored as (=) H=L♯L where L is (=) a linear operator with pseudo-adjoint L♯. In particular, this factorization applies to (e&eb) PT-symmetric and Hermitian Hamiltonians. The non-degenerate two-level systems provide (eb) a class of Hamiltonians that are pseudo-Hermitian. Authors demonstrate the implications of general results for (e&eb) this class in some detail. Copyright © 2002 Elsevier Science B.V. All rights reserved. Nuclear Physics B Volume 640, Issue 3, 30 September 2002, Pages 419–434Pseudo-supersymmetric quantum mechanics and isospectral pseudo-Hermitian Hamiltonians Ali Mostafazadeh STOP Michele Caselle, Alessandro Nada, Marco Panero present a high-precision lattice calculation of (e) the equation of state in the confining phase of (e) SU(2) Yang-Mills theory. Authors show that the results are described very well by (e) a gas of massive, non-interacting glueballs, provided (e) one assumes an exponentially growing Hagedorn spectrum. The latter can be derived within (eb) an effective bosonic closed-string model, leading to (eb) a parameter-free theoretical prediction, which is in perfect agreement with (=) proposed lattice results. Furthermore, when applied to (e&eb) SU (3) Yang-Mills theory, this effective model accurately describes (eb) the lattice results reported by Bors\'anyi et al. in JHEP 07 (2012) 056. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 07 (2015) 143 DOI: 10.1007/JHEP07 (2015)143 Cite as: arXiv:**

**1505.01106 [hep-lat] (or arXiv: 1505.01106v2 [hep-lat] for this version) Hagedorn spectrum and thermodynamics of SU (2) and SU (3) Yang-Mills theories Michele Caselle, Alessandro Nada, Marco Panero STOP B. de Wit J. Hoppe, H. Nicolai study the quantum-mechanical properties of (e) a supermembrane and examine the nature of (e) its ground state. A supersymmetric gauge theory of (e) area-preserving transformations provides (eb) a convenient framework for this study. The supermembrane can be viewed as (=) a limiting case of a class of models in supersymmetric quantum mechanics. Its mass does not depend on (=) the zero modes and vanishes only if (e) the wave function is (=) a singlet under (e&eb) supersymmetry transformations of (e) the nonzero modes. Authors exhibit (eb) the complexity of (e) the supermembrane ground state and (e&eb) examine various truncations of these models. None of these truncations has massless states. Copyright © 1988 Published by Elsevier B.V. Nuclear Physics B Volume 305, Issue 4, 28 December 1988: Pages 545-581: On the quantum mechanics of supermembranes Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. B. de Wit J. Hoppe, H. Nicolai STOP In this note we explain how world-volume geometries of (e) D-branes can be reconstructed within (eb) the microscopic framework where (e) D-branes are described through (e&eb) boundary conformal field theory. Authors extract the (non-commutative) world-volume algebras from (e) the operator product expansions of (e) open string vertex operators. For branes in a flat background with (e&eb) constant non-vanishing B-field, the operator products are computed perturbatively to (e) all orders in the field strength. Resulting series coincides with (=) Kontsevich's presentation of the Moyal product. After extending these considerations to (e) fermionic fields authors conclude with some remarks on the generalization of approach to (e&eb) curved backgrounds. D-branes and deformation quantization Volker Schomerus1,2 Published 25 July 1999 Journal of High Energy Physics, Volume 1999, JHEP06(1999)STOP Authors employ the path integral approach developed in Gangopadhyay S. and Scholtz F. E., Phys. Rev.**

**Lett., 102 (2009) 241602 to discuss the (generalized) harmonic oscillator in (eb) a noncommutative plane. The action for (e) this system is derived in (e&eb) the coherent state basis with (e&eb) additional degrees of freedom. From this the action in the coherent state basis without (e) any additional degrees of freedom is obtained. This gives (eb) the ground-state spectrum of the system. Authors then employ the exact renormalization group approach to show (eb) that equivalence can be constructed between this (noncommutative) system and (e&eb) a commutative system Path integral action and exact renormalization group dualities for quantum systems in noncommutative plane Sunandan Gangopadhyay Published 25 June 2015 • Copyright © EPLA, 2015 STOP Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini present a lattice study of (e)the equation of state in Yang-Mills theory based on (e) the exceptional G(2) gauge group. As is well-known, at zero temperature this theory shares (e&eb) many qualitative features with real-world QCD, including (e) the absence of colored states in the spectrum and dynamical string breaking at (e) large distances. In agreement with previous works, authors show that at finite temperature this theory features (eb) a first-order deconfining phase transition, whose nature can be studied by (e) a semi-classical computation. Authors also show that the equilibrium thermodynamic observables in the deconfined phase bear striking quantitative similarities with (=) those found in SU (N) gauge theories: in particular, these quantities exhibit (eb) nearly perfect proportionality to the number of gluon degrees of freedom, and (e&eb) the trace anomaly reveals (eb) a characteristic quadratic dependence on (e&eb) the temperature, also observed in SU (N) Yang-Mills theories (both in four and in three spacetime dimensions). Authors compare our lattice data with (e&eb) analytical predictions from (e) effective models, and discuss their implications for (e) the deconfinement mechanism and (e&eb) high-temperature properties of (e) strongly interacting, non-supersymmetric gauge theories. Results give (eb) strong evidence for (e) the conjecture that the thermal deconfining transition is governed by (e&eb) a universal mechanism, common to (e) all simple gauge groups.**

**Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 03 (2015) 057 DOI: 10.1007/JHEP03 (2015)057 Report number: DESY-14-146; IFT-UAM/CSIC-14-076 Cite as: arXiv: 1409.8305 [hep-lat] (or arXiv: 1409.8305v2 [hep-lat] for this version Exceptional thermodynamics: The equation of state of G(2) gauge theory Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini STOP In U (1) lattice gauge theory in three spacetime dimensions, the problem of confinement can be studied analytically in (e&eb) a semi-classical approach, in terms of (e&eb) a gas of monopoles with (e&eb) Coulomb-like interactions. In addition, this theory can be mapped to (e&eb) a spin model via (e&eb) an exact duality transformation, which allows (eb) one to perform high-precision numerical studies of (e) the confining potential. Taking advantage of these properties, authors carried out an accurate investigation of (e) the effective string describing (eb) the low-energy properties of flux tubes in (eb) this confining gauge theory. Authors found striking deviations from (e) the expected Nambu-Goto-like behavior, and, for the first time, evidence for (e) contributions that can be described by a term proportional to the extrinsic curvature of the effective string worldsheet. Such term is allowed by (e) Lorentz invariance and its presence in the infrared regime of the U (1) model was indeed predicted by (e) Polyakov several years ago. Results show that this term scales as expected according to (e) Polyakov's solution, and becomes (=) the dominant contribution to (e) the effective string action in the continuum limit. Authors also demonstrate analytically that the corrections to (e) the confining potential induced by (e) the extrinsic curvature term can be related to (e&eb) the partition function of (e) the massive perturbation of (e) a c=1 bosonic conformal field theory. The implications of results for SU (N) Yang-Mills theories in (e&eb) three and in four spacetime dimensions are discussed. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th) Journal reference: JHEP 1501 (2015) 105 DOI: 10.1007/JHEP01 (2015)105 Report number: IFT-UAM/CSIC-14-052 Cite as: arXiv:**

**1406.5127 [hep-lat] (or arXiv: 1406.5127v2 [hep-lat] for this version) A different kind of string Michele Caselle, Marco Panero, Roberto Pellegrini, Davide Vadacchino STOP In the past ten years, the ideas of supersymmetry have been profitably applied to (e&eb) many nonrelativistic quantum mechanical problems. In particular, there is now a much deeper understanding of why certain potentials are analytically solvable and (e&eb) an array of powerful new approximation methods for (e) handling potentials which are not (e) exactly solvable. In this report, authors review the theoretical formulation of (e) supersymmetric quantum mechanics and discuss many applications. Exactly solvable potentials can be understood in terms of (e&eb) a few basic ideas which include (e) supersymmetric partner potentials (e&eb) shape invariance and (e&eb) operator transformations. Familiar solvable potentials all have (e) the property of shape invariance. Authors describe new exactly solvable shape invariant potentials which include (e) the recently discovered self-similar potentials as (=) a special case. The connection between inverse scattering,(e&eb) isospectral potentials and (e&eb) supersymmetric quantum mechanics is discussed and multisoliton solutions of (e) the K dV equation are constructed. Approximation methods are also discussed within (eb) the framework of supersymmetric quantum mechanics and in particular it is shown that a supersymmetry inspired WKB approximation is (=) exact for a class of shape invariant potentials. Supersymmetry ideas give (eb) particularly nice results for the tunneling rate in (eb) a double well potential and for improving (eb+) large N expansions. Authors also discuss the problem of (e) a charged Dirac particle in (eb) an external magnetic field and other potentials in terms of (e&eb) supersymmetric quantum mechanics. Finally, authors discuss structures more general than supersymmetric quantum mechanics such as(=) parasupersymmetric quantum mechanics in which there is symmetry between a boson and (e&eb) a para-fermion of order p. Copyright © 1995 Published by Elsevier B.V. Physics Reports Volume 251, Issues 5–6, January 1995, Pages 267–385Supersymmetry and quantum mechanics Fred Coopera, Avinash Khareb, Uday Sukhatmec STOP The free energy of (e) the maximally**

**supersymmetric SU (N) gauge theory at temperature T is expected to scale, in the large-N limit, as (=) N2T4 times a function of the 't Hooft coupling, f (gYM2N). In the strong coupling limit the free energy has been deduced from (e) the near-extremal 3-brane geometry, and its normalization has turned out to be Full-size image (<1 K) times that found in (eb) the weak coupling limit. In this paper authors calculate the leading correction to (e) this result in inverse powers of (e) the coupling, which originates from (e) the R4 terms in the tree-level effective action of (e) type IIB string theory. The correction to Full-size image (<1 K) is positive and of (e) order Full-size image (<1 K). Thus, f (gYM2N) increases as (e&eb) the 't Hooft coupling is decreased, in accordance with (e&eb) the expectation that is should be approaching 1 in (eb) the weak coupling limit. Authors also discuss similar corrections for (e) other conformal theories describing (eb) coincident branes. In particular, authors suggest (eb) that the coupling independence of (e) the near extremal entropy for (e) D1-branes bound to (e) D5-branes is related to (e&eb) the vanishing of (e) the Weyl tensor of AdS3 × S3. Copyright © 1998 Published by Elsevier B.V. Nuclear Physics B Volume 534, Issues 1–2, 23 November 1998, Pages 202–222Coupling constant dependence in the thermodynamics of N = 4 supersymmetric Yang-Mills theory Steven S. Gubsera, Igor R. Klebanovb, 1, Arkady A. Tseytlinc, 2 STOP A new approach to (e) N=2 supersymmetry based on (e) the concept of harmonic superspace is proposed and is used to give (eb) an unconstrained superfield geometric description of N=2 super Yang-Mills and supergravity theories as (=) well as of matter N=2 hypermultiplets. The harmonic N=2 superspace has (e) an independent coordinates, in addition to the usual ones, the isospinor harmonics ui+or- on (e&eb) the sphere SU (2)/U (1). The role of ui+or- is to relate the SU (2) group realised on the component fields to (e&eb) a U (1) group acting on (e&eb) the relevant superfields. Their introduction makes it possible to SU (2)-covariantise (eb (t), e&eb) the notion of Grassmann analyticity. Crucial for the construction is the existence of (e) an analytic subspace of the general harmonic N=2 superspace. The hypermultiplet superfields and (e&eb) the true prepotentials (pre-**

**prepotentials) of N=2 super Yang-Mills and supergravity are (=) unconstrained superfunctions over (e&eb) this analytic subspace. The pre-prepotentials have (e) a clear geometric interpretation as gauge connections with (e&eb) respect to the internal SU (2)/U (1) directions. A radically new feature arises: the number of gauge and (e&eb) auxiliary degrees of freedom becomes (=) infinite while (e) the number of physical degrees of freedom remains (=) finite. Other new results are the massive N=2 Yang-Mills theory and (e&eb) various off-shell self-interactions of hypermultiplets. The propagators for matter and Yang-Mills superfields are given. Unconstrained N=2 matter, Yang-Mills and supergravity theories in harmonic superspace A Galperin, E Ivanov, S Kalitzin, V Ogievetsky and E Sokatchev STOP By assuming the existence of a novel multipronged string state for (e) D-particles interacting with (e&eb) D-brane intersections in type IIA string theory, authors are able to derive a quantum mechanical description of (e) supersymmetric Reissner-Nordstrom black holes. A supersymmetric index calculation provides (eb) evidence for (e) this conjecture. The quantum mechanical system becomes (=) two decoupled conformal quantum mechanical systems in (eb) the low energy limit. The conformal quantum mechanics has expected properties of (e) a dual description of string theory on (e&eb) AdS2×S2. Subjects: High Energy Physics - Theory (hep-th) Journal reference: JHEP 0001:006, 2000 DOI: 10.1088/1126-6708/2000/01/006 Report number: LBNL-43734 Cite as: arXiv:hep-th/9907101 (or arXiv:hep-th/9907101v5 for this version) Supersymmetric Quantum Mechanical Description of Four Dimensional Black Holes Julie D. BlumSTOP Using recently proposed soliton equations authors derive a basic identity for the scaling violation of (e) N=2 supersymmetric gauge theories Σiai∂F/∂ai−2F=8πib1u. Here F is (=) the prepotential, ai's are (=) the expectation values of (e) the scalar fields in (eb) the vector multiplet, u=1/2 Tr <ϕ2> and b1 is the coefficient of (e) the one-loop β-function. This equation holds in (eb) the Coulomb branch of all N=2 supersymmetric gauge theories coupled with (e&eb) massless matter. TOHRU EGUCHI and SUNG-KIL YANG, Mod Phys. Lett. A 11, 131 (1996).**

**DOI: 10.1142/S0217732396000151 PREPOTENTIALS OF N=2 SUPERSYMMETRIC GAUGE THEORIES AND SOLITON EQUATIONS STOP The Coulomb branch of N = 2 supersymmetric gauge theories in four dimensions is described in general by (e) an integrable Hamiltonian system in (eb) the holomorphic sense. A natural construction of such systems comes from (e) two-dimensional gauge theory and (e&eb) spectral curves. Starting from this point of view, authors propose an integrable system relevant to (e&eb) the N = 2 SU (n) gauge theory with (e&eb) a hypermultiplet in the adjoint representation, and offer (eb) much evidence that it is correct. The model has (e) a Full-size image (<1 K) S-duality group (with the central element −1 of Full-size image (<1 K) acting as (=) charge conjugation); Full-size image (<1 K) permutes the Higgs, confining, and oblique confining phases in the expected fashion. We also study more exotic phases. Copyright © 1996 Published by Elsevier B.V. Nuclear Physics B Volume 460, Issue 2, 5 February 1996, Pages 299–334Supersymmetric Yang-Mills theory and integrable systems Ron Donagia, 1, Edward Wittenb, 2 STOP A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 present a first step towards generalizing (eb) the work of Seiberg and Witten on N = 2 supersymmetric Yang-Mills theory to (e&eb) arbitrary gauge groups. Specifically, authors propose a particular sequence of (e) hyperelliptic genus n − 1 Riemann surfaces to underly (e&eb) the quantum moduli space of SU (n)N = 2 supersymmetric gauge theory. These curves have (e) an obvious generalization to (e) arbitrary simply laced gauge groups, which involves (e&eb) the A-D-E type simple singularities. To support proposal, authors argue that the monodromy in (e&eb) the semiclassical regime is (=) correctly reproduced. We also give some remarks on a possible relation to string theory. Copyright © 1995 Published by Elsevier B.V. Physics Letters B Volume 344, Issues 1–4, 26 January 1995, Pages 169–175Simple singularities and N = 2 supersymmetric Yang-Mills theory A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 STOP Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 review and extend earlier work that uses (e) the AdS/CFT correspondence to relate the**

**black-hole–black-string transition of gravitational theories on a circle to (e&eb) a phase transition in maximally supersymmetric (1 + 1)-dimensional SU(N) gauge theories at (eb) large N, again compactified on (e&eb) a circle. Authors perform gravity calculations to determine (eb) a likely phase diagram for (e) the strongly coupled gauge theory. Authors then directly study the phase structure of (e) the same gauge theory, now at (eb) weak't Hooft coupling. In the interesting temperature regime for (e) the phase transition, the (1 + 1)-dimensional theory reduces to (e&eb) a (0 + 1)-dimensional bosonic theory, which they solve using (e) Monte Carlo methods. Authors find strong evidence that the weakly coupled gauge theory also exhibits (eb) a black hole–black string-like phase transition in (eb) the large N limit. Authors demonstrate that a simple Landau–Ginzburg-like model describes (eb) the behaviour near the phase transition remarkably well. The weak coupling transition appears to be (=) close to the cusp between a first-order and a second-order transition. Black-hole–black-string phase transitions in thermal (1 + 1)-dimensional supersymmetric Yang–Mills theory on a circle Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 Published 29 October 2004 • 2004 IOP Publishing Ltd Classical and Quantum Gravity, Volume 21, Number 22STOP Monitoring the orbits of stars around (e&eb) Sgr A* offers(eb)the possibility of detecting (e&eb) the precession of (e) their orbital planes due to (e) frame dragging, of (e) measuring the spin and quadrupole moment of (e) the black hole, and of (e)testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and (e&eb) tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are (=), in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for which an optical interferometer, such as GRAVITY, will detect orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to (e&eb) the timescale of precession due to (e) the quadrupole moment of (e) the black hole. As a result, this non-conservative effect is (=) a**

**potential source of systematic uncertainty in (eb) testing the no-hair theorem with (e&eb) stellar orbits. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI:10.1088/0004-637X/777/1/57 Cite as:arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this version) Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP Authors perform a phase space analysis of (e) a generalized modified gravity theory with (e&eb) non-minimal coupling between geometry and (e&eb) matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that (eb) in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, Authors prove that there is (=) an attractor solution to describe (eb) the late time accelerating universe when (e) the modified gravity is chosen in (eb) a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and supernova type Ia data sets, Authors find that (eb) the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as:arXiv: 1512.09281 [gr-qc] (or arXiv: 1512.09281v1 [gr-qc] for this version) Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang study the P−V criticality and phase transition in (eb) the extended phase space of (e) charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of the system and (e&eb) its conjugate quantity is (=) the thermodynamic volume of (e) the black hole. The black**

**holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb)no P−V criticality and phase transition appear, while for the black holes with a spherical horizon, even when (e) the charge of the black hole is absent, the P−V criticality and the small black hole/large black hole phase transition (eb) will appear, but it happens only in (eb) d=5 dimensions; when (e) the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to (e) have the P−V criticality and the small black hole/large black hole phase transition, there exists (eb) an upper bound for (e) the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the Gauss-Bonnet coefficient and Q is the charge of the black hole. Authors calculate the critical exponents at the critical point and find that for all cases, they are (=) the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09 (2013)005 Cite as: arXiv: 1306.6233 [gr-qc] (or arXiv: 1306.6233v4 [gr-qc] for this version) P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP paper studies the whole process of (e) gravitational collapse and (e&eb) accretion of (e) a massless scalar field in (eb) asymptotically flat spacetime. Two kinds of initial configurations are considered. One is the initial data without (e) black hole; the other contains (e) a black hole. Under suitable initial conditions, Authors find that multi-horizon will appear which mean (eb) that the initial black hole formed by (e) gravitational collapse or existing at the beginning will instantly expand and suddenly grow (e&eb, eb), rather than grows gradually in the accretion process. A new type of critical behavior is found around (e&eb) the instant expansion. The numerical computation shows (eb) that the critical exponents are (=) universal. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Computational Physics (physics. Comp-ph) Cite as: arXiv: 1604.03363 [gr-qc] (or arXiv: 1604.03363v1 [gr-qc] for this version)**

**Multi-horizon and Critical Behavior in Gravitational Collapse of Massless Scalar Zhoujian Cao, Rong-Gen CAI, Run-Qiu Yang STOP In this paper, Authors present the results on the gravitational collapse of (e) charged massless scalar field in asymptotically flat spacetime with(e&eb) a perfectly reflecting wall. Differing from previous works, Authors study the system in (eb) the double null coordinates, by which we could simulate (e&eb) the system until the black hole forms with (e&eb) higher precision but less performance time. Authors investigate the influence of charge on (e&eb) the black hole formation and the scaling behavior near the critical collapses. The gapless and gapped critical behaviors for (e) black hole mass and charge are studied numerically. Authors find that they satisfy (eb) the scaling laws for critical gravitational collapse but the gapped critical behavior is different from (e) its AdS counterpart. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1602.00112 [gr-qc] (or arXiv: 1602.00112v1 [gr-qc] for this version) Multiple critical gravitational collapse of charged scalar with reflecting wall Rong-Gen CAI, Run-Qiu Yang STOP In this paper Authors investigate the (2+1)-dimensional topological non-Fermi liquid in (eb) strongly correlated electron system, which has (e) a holographic dual description by (e) Einstein gravity in (3+1)-dimensional anti-de Sitter (AdS) space-time. In a dyonic Reissner-Nordstrom black hole background, Authors consider a Dirac fermion coupled to (e&eb) the background U (1) gauge theory and an (e&eb) intrinsic chiral gauge field bM induced by (e) chiral anomaly. UV retarded Green's function of (e) the charged fermion in (eb) the UV boundary from (e) AdS4 gravity is calculated, by imposing (e&eb) in-falling wave condition at the horizon. Authors also obtain IR correlation function of (e) the charged fermion at (eb) the IR boundary arising from (e) the near horizon geometry of (e) the topological black hole with (e&eb) index k=0,±1. By using (e) the UV retarded Green's function and IR correlation function Authors analyze the low frequency behavior of (e) the topological non-Fermi liquid at zero and (e&eb) finite temperatures, especially the relevant non-Fermi liquid behavior near (eb) the quantum critical region. In addition, Authors**

**find that when k=±1, the effective mass of (e) bulk Dirac fermion at IR fixed point is (=) topologically quantized, consequently the Fermi momentum presents (eb) a discrete sparse distributed pattern, in contrary to (e&eb) the Ricci flat case with k=0. As demonstrated examples, Authors calculate the spectral functions of (e) topological non-Fermi liquid due to (e) anomalous Hall effect (AHE) (|b→|≠0,b 0=0) and chiral magnetic effect (CME) (|b→|=0,b0≠0), which can be tested in topological materials such as (=) topological insulator and (e&eb) topological metal with (e&eb) strong couplings. Subjects: High Energy Physics - Theory (hep-th); Strongly Correlated Electrons (cond-mat.str-el); General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv: 1601.03865 [hep-th] (or arXiv: 1601.03865v1 [hep-th] for this version) Topological Non-Fermi Liquid Rong-Gen Cai, Yong-Hui Qi, Yue-Liang Wu, Yun-Long Zhang STOP The unusual magnetic materials are significant in both science and technology. However, because of (e) the strongly correlated effects, it is difficult to understand their novel properties from (e) theoretical aspects. Holographic duality offers (eb) a new approach to understanding such systems from (e) gravity side. This paper will give a brief review of Authors recent works on the applications of holographic duality in understanding (e&eb) unusual magnetic materials. Some quantitative compare between holographic results and (e&eb) experimental data will be shown and some predictions from (e) holographic duality models will be discussed. Comments: For Second LeCosPA Symposium at Taiwan National University, Taipai, Dec. 14-18, 2015 Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1601.02936 [hep-th] (or arXiv:1601.02936v1 [hep-th] for this version) Understanding strongly coupling magnetism from holographic duality Rong-Gen Cai, Run-Qiu Yang STOP It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a**

**Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3**

**[astro-ph.CO] for this version) Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang**

# Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

## Module Number Six

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-465) 1990. We demonstrate its applicability cosmology/ quantum computation economics etal. For the first time, physicists couple distant nuclear spins using a single electron July 11, 2016. Some sentences have been deleted due to spatial restrictions. Please excuse me. For the first time, researchers at the University of Basel in Switzerland have coupled the nuclear spins of distant atoms using (e) just a single electron. Three research groups took part in this complex experiment, the results of which have now been published in the journal Nature Nanotechnology. In most materials, the nuclear spins of (e) neighboring atoms have (e) only a very weak effect on one another, as (=) the tiny nuclei are located deep within (e&eb) the atoms. This is not the case in metals, however, where (e) some of the electrons are able to move freely. The electron spins are able to couple nuclear spins at (eb) relatively large distances from one another. Named after four physicists, this RKKY interaction was discovered in the 1950s. Nuclear spins linked by (e&eb) an individual electron spin Now, an experiment by researchers from the Department of Physics at the University of Basel has for the first time succeeded in demonstrating this mechanism with (e&eb) a single electron, describing it using (e) quantum theory. The team led by Prof. Richard Warburton trapped a single electron inside (eb) a quantum dot. With use of a method developed in Basel to measure (eb) nuclear spin resonance, they showed that the electron coupled (e&eb) the spins of nuclei up to five nanometers apart - a huge distance in the world of (e) nuclear spins. The results are particularly relevant to (e&eb) the development of spin qubits: these seek to use (e) electron spins to carry (e&eb) information, but the interaction with (e&eb) the nuclei limits (e) the stability of the quantum information. Concentrated blast of physics More information: Role of the electron spin in determining the coherence of the nuclear spins in a quantum dot, Nature Nanotechnology, DOI: 10.1038/nnano.2016.114 Journal reference: Nature Nanotechnology search and more info website Provided by: University of Basel Read more at: http://phys.org/news/2016-07-physicists-couple-distant-nuclear-electron.html#jCp A new spin on quantum computing: Scientists train electrons with microwaves February 15, 2016 A silicon sample doped with (e&eb) bismuth atoms (left image) that is (=) just 150 nanometers thick is fitted with (e&eb) a superconducting resonator that includes (e) a capacitor (black, in left image; light gray in center image) and an inductive wire (red line … In what may provide (eb) a potential path to processing information in (eb) a quantum computer, researchers have switched an intrinsic property of electrons from an excited state to (e&eb) a relaxed state on demand using (e) a device that served as (=) a microwave "tuning fork." The team's findings could also lead to (eb) enhancements in magnetic resonance techniques, which are (=) widely used to explore (e) the structure of materials and biomolecules, and for medical imaging. The**

**international research team, which included scientists at the Department of Energy's Lawrence Berkeley National Laboratory (Berkeley Lab), demonstrated how to dramatically increase the coupling of (e) microwaves in a specially designed (e&eb) superconducting cavity to (e) a fundamental electron property called spin—which, like a coin, can be flipped. By zapping an exotic silicon material developed at Berkeley Lab with the microwaves, they found that they could rapidly change (e&eb) the electron spins from (e) an excited state to a relaxed, ground state by (e) causing (eb) the electrons to emit (eb) some of their energy in the form of (e) microwave particles known as photons. Left on their own, the electron spins would be extremely unlikely to flip back to (e) a relaxed state and to also emit (eb) a photon - the natural rate for this light-emitting effect, known as the Purcell effect, is about once every 10,000 years. The experiment demonstrated an accelerated, controllable relaxation of (e) electron spins and the release of (e) a microwave photon in about 1 second, said Thomas Schenkel, a physicist in Berkeley Lab's Accelerator Technology and Applied Physics Division who led the design and development of the silicon-bismuth sample used in the experiment. A new spin on quantum computing: Scientists train electrons with microwaves Like the classical tuning fork slung from guitar strings in this photo, scientists used (e) a device that worked like a microwave tuning fork to switch (e&eb) a fundamental property of electrons on (e&eb) demand in a silicon-bismuth sample developed at Berkeley Lab. Credit: Flickr/Pierre Guinoiseau "It's like a juggler who throws the balls up, and the balls come down 1,000 times faster than normal, and they also emit (eb) a microwave flash as they drop," he said. The results were published online Feb. 15 in the journal Nature. "Our results are highly significant for (e&eb) quantum information processing," said Patrice Bertet, a quantum electronics scientist at the French Atomic Energy Commission (CEA) who led the experiment. "Indeed, they are a first step toward (e&eb) the strong coupling of (e) individual electron spins to microwave photons, which could form (eb) the basis of new spin-based quantum computer architecture." John Morton, a professor at the London Center for Nanotechnology and co-author of the study, said, "Our ultimate aim is to find a link between quantum information that is fixed and (e&eb) quantum information that can be transported by photons." In today's computers, information is stored as (=) individual bits, and each bit can either be (=) a one or a zero. Quantum computers, though, could conceivably be (=) exponentially more powerful than modern computers because (e) they would use a different kind of bit, called a qubit, which because of (e) the weird ways of quantum mechanics can simultaneously behave (e&eb) as both a one and a zero. A coupled array of qubits would allow (eb) a quantum computer to perform (e&eb) many, many calculations at (eb) the same time, and electron spins are (=) candidates for qubits in (eb) a quantum computer. The latest study shows how the microwave photons could work in concert with (e&eb) the spins of electrons to move information in a new type of computer. "What we need now is ways to wire up these systems—to couple (e&eb) these spins together," Morton said. "We need to make coupled qubits that can perform (e&eb) computations." In the experiment, conducted at CEA in France, a small sample of a highly purified form of (e) silicon was implanted with (e&eb) a matrix of bismuth atoms, and a superconducting aluminum circuit was deposited on (e&eb) top to create (eb) a high-quality resonant cavity that allowed (eb) precise tuning of the microwaves. The electron spins of the bismuth atoms were then flipped into (e&eb) the excited, "spin-up" state. The microwave cavity was then tuned (e&eb) like a musical tuning fork, to (e) a particular resonance that coaxed the spins into emitting (eb) a photon as they flipped back to (e) a relaxed state. The cavity boosted the number of states into (e&eb) which a photon can be (=) emitted, which greatly increased (eb) the decay rate for (e) the electron spins in (eb) a controllable way. The technique is much like buying more lottery tickets to increase (eb) your chances of winning, Morton said. A new spin on quantum computing: Scientists train electrons with microwaves From left are, Berkeley Lab scientists Thomas Schenkel, Qing Ji and Peter Seidl at the NDCX-II (Neutralized Drift Compression Experiment II), which produces (eb) powerful ion beams. Researchers are exploring how to use NDCX-II to process exotic … The large bismuth atoms embedded in (eb) the silicon sample provided (eb) the electrons with (e&eb) unique spin properties that enabled (eb) the experiment. Schenkel said that implanting the bismuth atoms into (e&eb) the delicate silicon framework, a**

**process known as "doping," was "like squeezing bowling balls into (e&eb) a lattice of ping-pong balls." "We did a new trick with silicon. People wouldn't expect you could squeeze (e) anything new out of (e) silicon," Schenkel said. "Now we're looking into further improving (eb) bismuth-doped silicon and into (e&eb) tailoring the spin properties of (e) other materials, and using (e) this experimental technique for these materials." To enhance the performance of materials used (e) in future experiments, Schenkel said it will be necessary to improve (eb) the doping process so it is less damaging to (e) the silicon lattice. Also, the implantation process could be designed to produce (eb) regularly spaced arrays of (e) individual electron spins that would be more useful for (e) quantum computing than a concentrated ensemble of (e) electron spins. "We are now doing experiments on processing (e&eb) this and other materials at higher temperature and pressure with nanosecond ion pulses at NDCX-II, one of the accelerators here at Berkeley Lab," Schenkel said. "There are indications that it will improve (eb) the overall spin quality." Researchers said the latest research could potentially prove useful in boosting (eb) the sensitivity of scientific techniques like nuclear magnetic resonance spectroscopy and dynamic nuclear polarization, useful for a range of experiments, and could also shorten (e) experimental times by manipulating (e&eb) spin properties. "You need a way to reset (e&eb) spins—the ability to cause (eb) them to relax on demand to improve (eb) the rate at which you can repeat an experiment," Morton said. Bertet said it may be possible to further accelerate (eb) the electron-flipping behavior to (e) below 1 millisecond, compared to (e) the 1-second rate in the latest results. "This will then open the way to many new applications," he said. Explore further: Bismuth provides perfect dance partners for quantum computing qubits More information: A. Bienfait et al. Controlling spin relaxation with a cavity, Nature (2016). DOI: 10.1038/nature16944 Nature | Letter Controlling spin relaxation with a cavity A. Bienfait, J. J. Pla, Y. Kubo,X. Zhou, M. Stern, C. C. Lo, C. D. Weis, T. Schenkel, D. Vion,D. Esteve, J. J. L. Morton& P. Berte 531, 74–77 (03 March 2016) doi:10.1038/nature16944 Spontaneous emission of radiation is (=) one of the fundamental mechanisms by which an excited quantum system returns to (e) equilibrium. For spins, however, spontaneous emission is (=) generally negligible compared to (e&eb) other non-radiative relaxation processes because of (e) the weak coupling between the magnetic dipole and (e&eb) the electromagnetic field. In 1946, Purcell realized that the rate of spontaneous emission can be greatly enhanced by (e) placing the quantum system in a resonant cavity. This effect has since been used extensively to control (e&eb) the lifetime of atoms and semiconducting heterostructures coupled to (e&eb) microwave2 or optical3, 4 cavities, and is essential for (e) the realization of high-efficiency single-photon sources5. Here authors report the application of this idea to (e&eb) spins in solids. By coupling donor spins in silicon to (e&eb) a superconducting microwave cavity with (e&eb) a high quality factor and a small mode volume, authors reach the regime in which spontaneous emission constitutes (e) the dominant mechanism of spin relaxation. The relaxation rate is increased by (e) three orders of magnitude as the spins are tuned to (e) the cavity resonance, demonstrating that energy relaxation can be controlled (e&eb) on demand. Results provide a general way to initialize (eb) spin systems into their ground state and therefore have applications in (eb) magnetic resonance and (e&eb) quantum information processing6. They also demonstrate that the coupling between the magnetic dipole of a spin and (e&eb) the electromagnetic field can be enhanced up to (e) the point at which quantum fluctuations have (e) a marked effect on (e&eb) the spin dynamics; as such, they represent (eb) an important step towards the coherent magnetic coupling of (e) individual spins to (e) microwave photons Scientists obtain evidence of many-body localization in a closed quantum system July 8, 2016 by Olivia Meyer-Streng Science Cover Vol 352, Issue 6293, 24 June 2016, to highlight the research published in this issue (see original publication). The cover shows an artist's view of the experiment, illustrating the persistence of the density step of (e) an atomic distribution in an optical lattice. During equilibration ordinary many-body systems lose (e) all information about (e&eb) the initial state. Every morning authors experience an example for (e) this behaviour. Milk poured into a cup of coffee mixes (e&eb) perfectly and after some time it is impossible to say how exactly the two fluids were (=) put together. The same behaviour holds for (e) almost all quantum systems. However,**

**recently a new phenomenon called "many-body localization" has been predicted theoretically, which allows (eb) well insulated quantum systems to (e) preserve memory of the initial state forever. Now a team of scientists around Dr. Christian Groß and Professor Immanuel Bloch (Director at MPQ and Chair of Quantum Optics at LMU Munich), in cooperation with David Huse (Princeton University), has obtained evidence of such a behaviour in a two-dimensional quantum system of cold rubidium atoms trapped in (e&eb) an optical lattice. The scientists observed that – beyond a certain degree of disorder imprinted on (eb) the particle ensemble in the beginning – the system would relax into (e&eb) a steady state still containing (e) detailed microscopic information about its past. "We were able to observe the transition from a thermalized state into (e&eb) a many-body localized phase", Christian Groß points out. "It is the first observation of that kind in a regime that is not accessible with (e&eb) state-of-the-art simulations on classical computers." The experiment is not only of fundamental interest; the results might also lead to (eb) new ways for storing quantum information. Motivated by the foundational problem of how interacting particles behave (e&eb) in a disordered system, in the 1950s the American physicist Philip Warren Anderson discovered the famous localization phenomenon for (e) non-interacting particles, now called "Anderson localization". Here, disorder prevents (e) the particles to move and consequently all transport is (=) stopped. But what happens when disorder comes together with (e&eb) interactions? Will interactions lead to (eb) transport and thermalization, or will the localization persist even at (eb) high energies? So far, there is no theoretical model that faithfully predicts (eb) the evolution of a closed quantum system in more than one dimension under (e&eb) these conditions, although, the possibility for (e)localization has been theoretically suggested. In order to investigate these questions experimentally, strict requirements on controllability and perfect isolation of (e) the system have to be fulfilled. In the experiment described here, ultracold rubidium atoms are loaded into (eb) an optical lattice, a microscopic array of light traps formed by interfering (e&eb) laser beams. The disorder is created by projecting (e&eb) a computer-generated random light pattern onto (e&eb) the optical lattice. As a result of this "disorder", the depth of each microscopic trap in the lattice potential varies from site to (e) site. In fact, Prof. Bloch's group has advanced the experimental tools to such perfection that they can steer the position of (e) the atoms in the artificial light crystal and (e&eb) the interaction between (e&eb) them almost at will. With a high resolution microscope that detects the emitted fluorescence light, the position of each atom can be observed with (e&eb) high precision. Additionally, the atomic density distribution of the initial state can be (=) controlled, and for variable time intervals the evolution of the distribution can be (=) measured with high precision. These tools at hand, non-thermalizing behaviour can be probed in (eb) a conception ally very simple way. Any thermalized state of a closed system reflects (e&eb) the symmetry of its container – like, for example, water spilled into (e&eb) a round bowl would immediately cover (e) the whole bottom. In analogy, the scientists generate (eb) a density step in the initially prepared sample by (e) blowing away half of the atomic distribution with laser radiation. Then they watch how the remaining particles migrate into (e&eb) the empty half. For small disorders, the initially prepared density step is smeared out (e) fast and the initially empty and filled halves become (=) indistinguishable. However, when the measurement is repeated for (e) strong disorder, traces of the initial state remain and the system does not (e) relax to a thermal state even for (e) very long times. "We observe a fairly sharp onset of non-thermalizing behaviour above (e) a critical value", says Christian Groß. "This absence of thermalization is remarkable because (e) thermalization persists in a system of interacting particles, even at (eb) the high energies probed in the experiment." The scientists interpret their observation as (=) the onset of many-body localization in (eb) the atomic system. This is of fundamental interest because it means (eb) the breakdown of equilibrium statistical mechanics. On the other hand, the persistence of initial state information could be used as (=) a source for quantum information technologies. "It should also be emphasized that we obtain these results for a system size that is far beyond (e) numerically accessible scales", says Jae-yoon Choi, postdoc at the experiment. More information: J.-y. Choi et al. exploring the many-body localization transition in two dimensions, Science (2016). DOI: 10.1126/science.aaf8834 Journal reference: Science search and more**

**info website Provided by: Max Planck Society search and more info website Bosons refusing to thermalize in 2D Messy, interacting quantum-mechanical systems are difficult to analyze theoretically. In a single spatial dimension, the calculations are (=) still tractable, and experiments have recently confirmed (eb) the prediction that sufficiently strong disorder can disrupt (e&eb) the transport of interacting particles. In two dimensions, however, the theoretical blueprint is missing. Choi et al. used (e) single-site imaging of cold 87Rb atoms in (eb) an optical lattice to show (eb) that similar localization occurs in (eb) two-dimensional (2D) systems. The study highlights the power of quantum simulation to solve problems that are currently inaccessible to classical computing techniques. Science, this issue p. 1547 Abstract A fundamental assumption in statistical physics is that generic closed quantum many-body systems thermalize under (e&eb) their own dynamics. Recently, the emergence of many-body localized systems has questioned (e) this concept and challenged our understanding of the connection between statistical physics and (e&eb) quantum mechanics. Here authors report on the observation of a many-body localization transition between thermal and (e&eb) localized phases for (e) bosons in a two-dimensional disordered optical lattice. With our single-site–resolved measurements, authors track the relaxation dynamics of (e) an initially prepared out-of-equilibrium density pattern and find strong evidence for (e) a diverging length scale when (e) approaching the localization transition. Experiments represent a demonstration and in-depth characterization of many-body localization in (eb) a regime not accessible with (e&eb) state-of-the-art simulations on classical computers. Pelican instrument provides crucial experimental evidence of unusual quantum state July 8, 2016 Artist's illustration. Credit: Australian Nuclear Science and Technology Organisation (ANSTO) Experiments using inelastic neutron scattering at the Australian Centre for Neutron Scattering have found indications (eb) of a possible new quantum spin state in a novel antiferromagnetic material barium ytterbium zinc oxide (Ba3Yb2Zn5O11) which provides (eb) both a challenge and validation of (e) the third law of thermodynamics. The quantum mechanical information gained from (e) the experiment is crucial not only in understanding the quantum state of the specific material but also provides (eb) better understanding of high temperature superconductivity and ultimately offers (eb) guidance in the design of materials for (e) quantum computing. A research team led by Associate Professor Takatsugu Masuda at the Institute for Solid State Physics at the University of Tokyo with collaborators from ANSTO, Osaka University and Taiwan's National Cheng Kung University, have performed a successful experiment using Pelican, time-of-flight cold neutron spectrometer at ANSTO. The results have just been published in Physical Review B as a rapid Communication and an Editors' Suggestion. "The great success of this challenging experiment relied on (e) the excellent performance of the Pelican instrument, wonderful technical support at ANSTO and strong expertise in the relevant field from A/Prof Masuda's group," said Dr Dehong Yu and Dr Richard Mole, co-authors and Pelican instrument scientists. Pelican instrument provides crucial experimental evidence of unusual quantum state INS spectra measured at 1.5 K (a), 6 K (b), 12 K (c), and 40 K (d) using setup I. (e) INS spectrum measured at 1.5 K using the setup II. (f) Q dependence of the integrated intensity obtained from spectrum in (e). Red solid curve is the …….(deleted due to spatial constraints) Barium ytterbium zinc oxide (Ba3Yb2Zn5O11) is a novel antiferromagnetic model system having (e) a "breathing pyrochlore lattice" consisting of (e) arrays of alternating large and (e&eb) small corner-sharing tetrahedra of ytterbium atoms. Earlier research had suggested (eb) the material was (=) an isotropic spin half (S= ½) system. However, four low energy excitations measured on the sample with the Pelican instrument at a temperature of 1.5 K clearly disagree with (e) this designation and the associated theoretical model based on (e) the isotropic spins having simple interactions. "These observations suggested (eb) a much more complex spin state in the system," said Yu. "The experiment is interesting and important because as you start going very low in temperature, quantum mechanics takes over (e) system rather than thermodynamics, and you observe unusual properties," said Mole. "According to low temperature heat capacity measurements, the material does not (e) have zero entropy or a single ground state until you reach 63 milli Kelvin or -272.94 Celsius, before it starts obeying the third law." Pelican instrument provides crucial experimental evidence of (e) unusual quantum state**

**Breathing pyrochlore lattice formed by (e) A3+ ions as determined by Kimura, Nakatsuju and Kimura in previous work The third law of thermodynamics states that for pure crystals as the temperature of a system approaches absolute zero K, (eb) its entropy or disorder approaches zero. "The research clearly indicates (eb) that below this temperature the material may be in (eb) an exotic state, such as (=) a liquid-like spin state, in which the arrangement of spins behaves (e&eb) like a liquid. In this exotic state, novel phenomena may show up and be detectable; for example, charge and spin separation of electrons may be (=) possible," said Yu. The researchers suggest that further investigations, experimentally and theoretically, are required to ultimately determine the true identity of (e) the ground state of the material. "Discovering the existence of regular tetrahedral spin made it well worth coming all the way to Australia to make use of this neutron experimental apparatus," said Associate Professor Masuda. Many Japanese researchers are using ANSTO's neutron scattering instruments through the peer review system, while their nuclear facilities are offline. The outcome of this particular experiment is also part of the collaboration agreement between ANSTO and the University of Tokyo. More information: T. Haku et al. Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet, Physical Review B (2016). DOI: 10.1103/PhysRevB.93.220407 Journal reference: Physical Review B search and more info website: provided by: Australian Nuclear Science and Technology Organisation (ANSTO) search and more info website Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 T. Haku, K. Kimura, Y. Matsumoto, M. Soda, M. Sera, D. Yu, R. A. Mole, T. Takeuchi, S. Nakatsuji, Y. Kono, T. Sakakibara, L.-J. Chang, and T. Masuda Phys. Rev. B 93, 220407(R) – Published 24 June 2016 Authors study low-energy excitations in (eb) the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 by (e) a combination of inelastic neutron scattering (INS) and (e&eb) thermodynamical property measurements. The INS spectra are quantitatively explained by (e) spin-12 single-tetrahedron model having (e) XXZ anisotropy and Dzyaloshinskii-Moriya interactions. This model has a twofold degeneracy of (e) the lowest-energy state per tetrahedron and well reproduces (eb) the magnetization curve at (eb) 0.5 K and heat capacity above 1.5 K. At lower temperatures, however, authors observe a broad maximum in (eb) the heat capacity around 63 mK, demonstrating (eb) that a unique quantum ground state is selected due to (e) extra perturbations with (e&eb) an energy scale smaller than the instrumental resolution of INS. Revised 13 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.220407 © 2016 American Physical Society Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA (Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor. In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface. As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone. Such planets are all the "right" distance from (e) their star, and hopefully, also have water. But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone. For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement. But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway. But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain. It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of**

**reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life. Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface. The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113 Abstract The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question. Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry. This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry. Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments. Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy. The electronic and (e&eb) structural variability among them is extraordinary. The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations. The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface. The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers. Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp**

## 3. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_6}{dt} = -b_6 Q_6 + g_6(g_6)Q_6$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_6(0) = 0, \quad g_6(P_6) > 0, \quad P_6 \ 0, \qquad \lim_{P_6 \to \infty} g_6(P_6) = (K_6)_6 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_6 + (\rho_a)_6 + D_6(a_6, P_6, Q_6)\rho_6 + 0, \qquad 0 < a_6 \ \infty,$$

$$\rho_6(t, 0) = \int_0^\infty n_6(a_6, P_6, Q_6)\rho_6(t_6, a_6)da_6$$

Where $D_6$ is the mortality function and $n_6$ is the fertility function.

We assume following equations held

$$\left[\frac{dQ_6}{dt}\right]_6 = x_6\left[\frac{dQ_2}{dt}\right]_6 - x_2\left[\frac{dQ_6}{dt}\right]_6$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_6 - x_4\left[\frac{dQ_2}{dt}\right]_6$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_6 - x_6\left[\frac{dQ_3}{dt}\right]_6$$

Where $x_6, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_6$ and $g_6(P_6)Q_6$

Now when you add the term $(b_6Q_6)$ with $g_6(P_6)(Q_6)$ we are infect see is that $(b_6Q_6)$ is credited incrementally part of $g_6(P_6)(Q_6)$ or instant nearly, when subtracted from $\frac{dQ_6}{dt}$ , it means that $g_6(P_6)(Q_6)$ is slowly eaten by $\frac{dQ_6}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_6 = \mu_6(a_6) + (f_6)_6(P_6)_6 + (f_2)_6(Q)_6 \quad (2)$$

$$n_6 = \beta_6(a_6)h_6(Q_6), \quad (3)$$

Where $\mu_6$the age- dependent debit rate whit is $\mu_6(a_6)$ 0 $\forall a > 0$; $(f_6)_6$ is the density- dependent death rate with $f_6(0) = 0$, $(f_6^6)_6(P_6)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_6)_6(P_6) = (k_6)_6$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_6(0) = 0, (f_2^6)_6(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_6(Q_6) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 6, h^6(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_6 + (\rho_a)_6 = -\mu_6(a_6) + (f_6)_6P_6) + \big((f_2)_6(Q_6)\big)(\rho_6)_6 \quad (4)$$

$$\rho_6(t_6, 0) = h_6(Q_6)\int_0^\infty \beta_6(a_6)\rho_6(t, a_6)da, \quad (5)$$

$$\frac{dQ_6}{dt} = -b_6Q_6 + g_6(P_6)Q_6 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_6$ and $Q_6$ of (4-6) satisfy

$$0 \ \rho_6(t, a_6) \ \rho_6^*(t, a_6),$$

$$0 \leq Q_6(t) \leq Q_6^*(t),$$

Where $\rho_6^*(t, a_6)$ and $Q_6^*(t)$, are solutions of

$$(\rho_t^*)_6 + (\rho_a^*)_6 = -(\mu_6(a_6) + (f_6)_6(P_6^*) + (f_2)_6 Q_6^*))\,\rho_6^*, \quad (7)$$

$$\rho_6^*(t, 0) = \int_0^\infty \beta_6(a_6)\,\rho_6^*(t, a_6)da_6 \quad (8)$$

$$\frac{dQ_6^*}{dt} = -b_6\,Q_6^* + g_6(P_6^*)Q_6^*, \quad (9)$$

With $P_6^*(t) = \int_0^\infty \rho_6^*(t, a_6)da_6$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_6^*(t)}{dt} = P_6^*(t)\big(r_6 - (f_6)(P_6^*) - (f_2)_6(Q^*)\big), \quad (10)$$

$$\frac{dQ_6^*(t)}{dt} = Q_6^*(t)(-b_6 + g_6(P_6^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_6(a)e^{-\int_0^a \mu_6(\tau)dt}\,e^{-ra}\,da_6 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_6^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(t)dt}\, da_6 > 1, \qquad (H1)$$

Which guarantees $r_6 > 0$.

*Remark.* If $\beta_6(a) = [\beta_0]_6 a_6 e^{-\alpha_6 a_6}, \mu_6(a_6) \equiv \mu_6$ then the solution of (12) is

$$r_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_6]_6(P_6) = 0, [f_2](Q_6) = \lambda_6 Q_6\ h_6(Q_6) = 1, \qquad and\ \ g_6(P_6) = c_6 P_6$$

$$[\rho_6]_6 + [\rho_a]_6 = -(\mu_6 + \lambda_6 Q_6)\rho_6$$

$$\rho_6(t,0) = \int_0^\infty [\beta_0]_6 a_6 e^{-\alpha_6 a_6} \rho_6(t, a_6)\, da_6$$

$$\frac{dQ_6}{dt} = -b_6 Q_6 + c_6 P_6 Q_6$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_6}{dt} = \omega_6 P_6 - \lambda_6 P_6 Q_6$$

$$\frac{dQ_6}{dt} = b_6 Q_6 - c_6 P_6 Q_6$$

Where $\omega_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$

## 4. THE EQUILIBRIUM $(\rho_0)_6(a_6), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_6(a_6), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_6 + (\rho_a)_6 = -\big(\mu_6(a_6) + (f_6)_6(P_6)\big)\rho,$$

$$\rho_6(t,0) = \int_0^\infty \beta_6(a_6)\rho_6(t,a_6)da_6 \qquad (13)$$

The solution of (63) has the form of

$$\rho_6(t,a_6) = Ce_6^{r(t-a_6)} - \int_0^a (\tau)d\tau - \int_0^6 (f_6)_6\big(P_6(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_6(t) = C_6 e^{r_6 t - \int_0^t (f_6)_6(P_6(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_6^{-6}(r)$ if $k_6 > r$. This is equivalent to

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{(k_6)_6 a_6} da_6 < 1.$$

In fact, if we define

$$R_6(s_6) = \int_0^\infty \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau)d\tau} e^{-s_6 a_6}\, da_6$$

$R(s)$ is a decreasing function. $R_6(r_6) = 1$ And $R_6(k_6)_6 < 1$ imply $r_6 < (k_6)_6$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_6(a_6), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{k_6 a_6} da_6$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_6 \equiv \int_0^\infty (\rho_0)_{66}(a_6)da_6$ *satisfies*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{-(f_6)_6(P_0)a_6} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_6(a_6), 0)$

Let $u_6(t, a_6) = \rho(t, a_6) - \rho_0(a_6)$ Then

$$N_6(t) \equiv \int_0^\infty u_6(t, a_6)da_6 = \int_0^\infty \rho_6(t, a_6)da_6 - \int_0^\infty (\rho_0)_6(a_6)da = P_6(t) - P_0)_6$$

And

$$(u_6)_6 + (u_a)_6 = (\rho_t)_6 + (\rho_a)_6 - (\rho_0^6)_6(a_6)$$

$$= -\mu_6(a_6) + (f_6)_6(P_6) + (f_2)_6(Q_6))\,(u_6 + (\rho_0)_6) + \mu_6(a_6)(\rho_0)_6 + (f_6)_6(P_0)_6(\rho_0)_6$$

$$= -\mu_6(a_6)u_6 - (f_6)_6(P_6) + (f_2)_6(Q_6))\,u_6 - (\rho_0)_6\,(f_6)_6(P_6) - (f_6)_6(P_0)_6$$

$$+ (f_2)_6(Q_6)), \qquad (15)$$

Where $(\rho_0^6)_6(a_6)$ denotes $d(\rho_0)_6(a_6)/da_6$

Expansion $(f_6)_6(P_6)$ and $(f_2)_6(Q_6)$ at $(P_0)_6$ and 0 yield respectively

$$(f_6)_6(P_6) = (f_6)_6((P_0)_6 + N_6)$$

$$= (f_6)_6(P_0)_6 + (f_6^6)_6(P_0)_6 N_6 + o_6(N_6), \qquad (16)$$

$$(f_2)_6(Q_6)$$

$$= (f_2^6)_6(0)Q + o_6(Q_6) \quad (17)$$

Write $f_6(P_0) = f_{60}, f_6'(P_0) = f_{60}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_6)_6 + (u_a)_6 \approx -(\mu_6(a_6) + (f_{60})_6)u - (f_{60}^6)_6(\rho_0)_6(N_6) - (f_{20}^6)_6(\rho_0)_6(Q_6)$$

$$u_6(t,0) = \rho_6(t,0)(\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a) + (\rho_0)_6(a_6)da_6 - (\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a_6)da_6 + (h_6(Q_6) - 1)(\rho_0)_6(0) \quad (18)$$

Expanding

$$h_6(Q_6) = 1 + (h^6)_6(0)Q_6 + 0(Q_6) \equiv 1 + h_0'Q_6 + 0Q_6$$

And substituting it into (18), we arrive at

$$u_6(t,0) \approx \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + (h_{a_6}')(\rho_0)_6Q_6$$

The linearization is the following :

$$[u_6] + [u_6] = -(\mu_6(a_6) + [t_{60}]_6u_6 - [\alpha_6]_6(a_6)N_6 - [\alpha_2]_6(a_6)Q_6$$

$$u_6(t,0) = \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + [\xi_6]Q_6 \quad (19)$$

$$\frac{dQ_6}{dt} = -b_6Q_6$$

where

$$[\alpha_6]_6(a_6) = [t_6^6 a]_6[p, o]_6(a_6)d2(a_6) = [f_{20}^6]_6(\rho\, o\,)_6(a_6)\xi = [\mathrm{h\ o}^6]_6[\rho \mathrm{o}]_6(\mathrm{o}) \text{ and } \mathrm{b}_6 = \mathrm{b}_6 - g_6(P_6)$$

If $g_6(P_0) > b$, clearly, $(\rho_0)_6(a_6), 0)$ is unstable and the predator population is persistent. Assume $g_6(P_0)_6 < b$.

Then, $Q_6(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_6 + u_a = -(\mu_6(a_6) + (f_{60})_6)u_6 - (\alpha_6)_6(a_6)N_6$$

$$= -(\mu_6(a_6) + f_{60})u_6 - \alpha_6(a_6)\int_0^\infty u_6(t, a_6)da_6$$

$$u_6(t, 0) = \int_0^\infty \beta_6\,(a_6)u_6(t, a_6)da_6$$

Using the method developed by Cushing [2, 3], we suppose

$$u_6(t, a_6) = \tilde{\mathrm{u}}_6(a_6)^{e^{\gamma(t-a_6)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_6}\, da_6$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_6)$ gives,

$$\tilde{\mathrm{u}}_6(a_6) = e^{-\int_0^a \mu_6(\tau_6)d\tau - (f_{60})_6 a}\left(\tilde{\mathrm{u}}_6(0) - w\int_0^a e^{\int_0^\tau \mu_6(s_6)ds + (f_{60})}\tau_6\,(\alpha_6)_6(\tau_6)e^{\gamma_6\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_6(a_6)e^{-\gamma_6 a_6 - (f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6)d\tau}\right)\tilde{\mathrm{u}}_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a(\alpha_6)_6(\tau_6)e^{-\gamma_6\tau_6}e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds}d\tau_6\,da_6$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}da\right)(\tilde{u}_6(0))_6$$

$$+\left(-6-\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds+y\tau}\alpha_6(\tau)d\tau\,da\right)w=0. \qquad (23)$$

Define $Y(\gamma,a)=e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(t_6)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_6(a_6)Y_6(\gamma_6,a_6)da_6\right)u_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6=0, \qquad (24)$$

$$\left(\int_0^{\infty}Y_6(\gamma_6,a_6)da_6\right)\tilde{u}_6(0)$$

$$+\left(-1-\int_0^{\infty}Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6\right)w=0.$$

Equqtion (20) has a solution of the form of $u_6(t,a_6)=u_6(a_6)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_6(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) Y^{-6}(\gamma_6, a_6) d\tau_6$. We obtain the characteristic equation

$$f_{60}^{l}{}_6 \left( \int_0^{\infty} Y_6(\gamma_6, a_6) da_6 \right) \left( \int_0^{\infty} \beta_6(a_6) Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$+ \left( 1 - \int_0^{\infty} \beta_6(a_6) Y_6(\gamma_6, a_6) da \right) \left( 1 + (f_{60}^{l})_6 \int_0^{\infty} Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$= 0. \qquad (25)$$

Since

$$(\rho_0)_6(a_6) = (\rho_0)_6(0) e^{-\int_0^a \mu_6(\tau_6) d\tau_6 - (f_{60})_6 a_6}$$

$$I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) e^{\gamma_6 \tau_6 + (f_{60})_6 \tau + \int_0^{\tau} \mu_6(s_6) ds_6} d\tau$$

$$= \frac{(\rho_0)_6(0)}{\gamma_6} (e^{\gamma_6 a_6} - 1)$$

Let

$$E_6(a_6) = e^{-(f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6) d\tau_6}$$

This relationship is cardinal one. It is useful in quantum computation and other black hole dynamics problems

**Quantum space and quantum time**

(111) **Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.**

(112) **In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.**

(113) **They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.**

(114)**When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.**

(115) **Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.**

**The Planck scale**

(116)**Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".**

(117)**Here ħ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.**

(118)**The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.**

(119)**Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)**

(120)**Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.**

## Physical applications

**Some phenomena where quantum gravity plays a major role are the following.**

**The microscopic structure of spacetime**

(121)**If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.**

(122) **Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length**

(123) **Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.**

(124) **A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.**

**Early cosmology**

(125)**According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.**

(126) **However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.**

(127) **A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.**

(128) **Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".**

(129)**One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.**

**Black holes**

(130)**Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.**

(131) **First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.**

(132) **Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.**

(133)**What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.**

(134) **Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.**

**Astrophysical effects**

(135)**Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.**

(136) **An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.**

(137)**The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.**

**Main approaches**

(138)**The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.**

(139) **Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.**

(140)**The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.**

**String theory**

(141)**String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.**

(142) **Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.**

(143)**String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.**

(144)**Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.**

(145)**However, when summed the entire series appears to be (=) divergent.**

(146)**In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.**

(147)**Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.**

(148)**Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.**

(149) **In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.**

**Loop quantum gravity and spin foams**

(150)**Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.**

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|---|
| (151)**The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity.** |
| (152)**Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime.** |
| (153)**The granular properties of space can be (=) explicitly computed.** |
| (154) **In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory.** |
| (155)**The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks.** |
| (156) **Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure.** |
| (157) **The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit.** |
| (158)**Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism.** |
| (159)**This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism.** |

**Noncommutative geometry**

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|---|
| (160)**Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime.** |
| (161) **Quantum mechanics teaches us that dynamical quantities are (=) noncommutative --meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed.** |
| (162)**It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry.** |
| (163) **A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces.** |
| (164)**Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces.** |
| (165)**This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity.** |

**Others**

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|---|
| (166) **Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays** |

(167)**. Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection.**

(168) **Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible.**

(169) **The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.**

**Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016**

(170)**These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field.**

(171)**The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets.**

(172) **Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state.**

(173) **Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands.**

(174)**The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI**

(175) **The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9].**

(176)**This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive.**

(177)**More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.**

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

(178)**Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.**

(179)**Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.**

(180) **The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.**

(181)**This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.**

(182)**The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.**

(183)**This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.**

(184) **For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.**

(185)**The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society**

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(186)**Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.**

(187)**We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.**

(188)**When this is true—as is often the case—the system serves as (=) a "bath" for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called thermalization.**

(189)**However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.**

(190)**Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1).**

(191)**Bringing them to thermal equilibrium requires (e) an external bath.**

(192) **But so far, there has been no rigorous proof of (e) the existence of such many-body localization. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption.**

(193)**The concept of localization was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations).**

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| (194)**The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century.** |
| (195)**But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information.** |
| (196)**Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment.** |
| (197)**For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood.** |
| (198) **It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7].** |
| (199) **Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature.** |
| (200) **One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.** |
| (201)**However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].** |
| (202)**Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.** |
| (203) **Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.** |
| (204)**As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked.** |
| (205)**His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice.** |
| (206) **An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site.** |
| (207)**This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself).** |
| (208)**This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition.** |
| (209)**However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime:** |

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| (210) **These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects.** |
| (211)**It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies).** |
| (212) **This is a realistic scenario for (e) all known systems.** |
| (213)**On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects.** |
| (214)**On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other.** |
| (215)**No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true.** |
| (216) **His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found.** |
| (217)**Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization?** |
| (218)**Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects?** |
| (219)**Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions.** |
| (220)**So in those systems, the quantum phase transition does indeed exist.** |
| (221)**But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension.** |
| (222)**Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters.** |

**Focus: Biological Cells Form Electric Circuits July 1, 2016• Physics 9, 75 Cells that are electrically active and that also produce light for easy voltage monitoring could lead to new studies of heart arrhythmias and possibly bio-computing. Figure caption expand figure H. M. McNamara et al., Phys. Rev. X (2016) Cell circuit. Waves of electrical activity propagate in a loop within a collection of cells engineered with electrical and optical properties (see video below).Some sentences, pictures and videos are deleted due to constraints on space.**

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| (223)**The human heartbeat is produced by (e) electrical pulses that propagate through (e&eb) cardiac tissue, causing (eb) rhythmic muscle contraction.** |
| (224)**Researchers have previously engineered cells to create (eb) an artificial tissue capable of producing (eb) coordinated electrical activity, and now a team has added the ability to monitor (e&eb) their electrical state by (e) detecting fluorescent emission.** |
| (225)**They have also fashioned the cells into (e&eb) "living circuits" that might act as (=) model systems for studying heart behavior.** |
| (226)**Constructing electrically active tissues from their basic building blocks by (e) genetically engineering cells allows (eb) researchers to study a "clean" system similar to (e) the more complicated cardiac tissue.** |

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| (227)**These systems let them investigate, for example, the electrical signals that lead to (eb) the potentially fatal irregular heartbeat called arrhythmia.** |
| (228)**Previously researchers have introduced (eb) ion channels—proteins that conduct ions through (e&eb) cell membranes—into human kidney cells to make (eb) electrically "excitable" tissues [1].** |
| (229)**Adam E. Cohen of Harvard University and his co-workers decided to add more ingredients to such cells, so that the electrical activity could be (=) light-activated and easily visualized.** |
| (230)**They incorporated four additional gene products into (e&eb) human embryonic kidney (HEK) cells.** |
| (231) **Two were ordinary ion channels; the other two were (=) a blue-light-activated ion channel (a light-sensitive switch) and (e&eb) a protein that fluoresces with (e&eb) an intensity that depends on (=) the voltage across the membrane in which it is embedded.** |
| (232)**Light activation of electrically excitable cells has been shown before [2]. But Cohen and colleagues have now included a convenient "readout" mechanism (the fluorescent protein) and have used patterning methods to shape (e&eb) the cell cultures into "tissue circuits."** |
| (233)**They grew rings of cells a few millimeters in diameter that could act as (e&eb) oscillators that support (eb) circulating waves of electrical excitation.** |
| (234)**These waves arise from (e) the collective behavior of the cells: a pulse produced by (e) one cell can trigger another from (e) its neighbors, eventually producing (eb) a coherent wave of excitation traveling through (e&eb) the tissue.** |
| (235) **It's this coordination, in a repeating cycle of excitation and quiescence, which gives rise to (eb) rhythmic activity in cardiac tissue.** |

**H. M. McNamara et al., Phys. Rev. X (2016)**

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| (236)**A pattern of blue light (not visible) illuminates these cells, defining (eb) a rectangular loop around which electrical waves can propagate.** |
| (237)**At about 0:07 the rectangle begins to shrink in (eb) the vertical direction, causing t(eb) he cycle time to gradually decrease. (In other experiments, the cells were grown in (eb) fixed, ring-shaped patterns.)** |
| (238)**The ring of engineered HEK cells acted like an electronic "ring oscillator," a device with a characteristic frequency to (e&eb) its circulating electrical signal.** |
| (239)**The electrical waves in a ring oscillator can propagate (e&eb) in either direction around the ring, enabling (eb) it to serve as a binary memory unit.** |
| (240) **Cohen and colleagues found that in their HEK tissues, these circulating waves remained (eb) stable for more than two hours, with (e&eb) the pulses completing a few circuits every second.** |
| (241)**The electrical signals were easily visualized as (=) waves of fluorescent emission.** |
| (242)**The researchers say that these tissues could be used to (e) understand and predict the behavior of real cardiac cells, especially arrhythmias.** |
| (243)**"In real heart, if you pace the tissue fast enough, *(eb) it starts skipping beats and eventually becomes (=) arrhythmic," says Cohen.** |
| (244) **This state is a well-known precursor to (e) sudden cardiac death and is associated with (e&eb) the breakup of the usual regular waves of electrical excitation.** |
| (245) **Some drugs also have the dangerous side effect of inducing (eb) arrhythmia, and Cohen says that the new system might offer a quick and convenient way of screening for such hazards. The work "is really very elegant and deserves attention," says physiologist Gil Bub of Oxford University in the** |

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| **UK. "It opens new research questions for people like me who are interested in wave dynamics in excitable systems."** |
| (246)**There might also be practical uses for information processing in living tissues. "By combining patterns of cell growth and (e&eb) patterns of ion channel expression, we think we can build diodes, amplifiers, and other simple logic elements," says Cohen.** |
| (247)**Andrew Adamatzky of the University of the West of England in Bristol, a specialist in "unconventional computing," foresees circuits that perform computations through (e&eb) the interactions of electrical waves. "This research might lead to fabrication of implantable bio-computers," he says. This research is published in Physical Review X. –Philip Ball Philip Ball is a freelance science writer in London and author of Invisible: The Dangerous Allure of the Unseen (University of Chicago Press, 2014).** |

**Optically Controlled Oscillators in an Engineered Bioelectric Tissue Harold M. McNamara, Hongkang Zhang, Christopher A. Werley, and Adam E. Cohen Phys. Rev. X 6, 031001 – Published 1 July 2016**

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| (248)**Complex electrical dynamics in excitable tissues occur throughout (e&eb) biology, but the roles of individual ion channels can be difficult to determine due to (e) the complex nonlinear interactions in (eb) native tissue.** |
| (249) **Here, authors ask whether they can engineer a tissue capable of (e) basic information storage and processing, where (e) all functional components are known and well understood.** |
| (250)**Authors develop a cell line with (e&eb) four transgenic components: two to enable (eb) collective propagation of electrical waves and (e&eb) two to enable optical perturbation and (e&eb) optical readout of membrane potential.** |
| (251)**They pattern the cell growth to define (eb) simple cellular ring oscillators that run stably for >2 h (∼104 cycles) and that can store data encoded in (eb) the direction of electrical circulation.** |
| (252) **Using (e) patterned optogenetic stimulation authors probe the biophysical attributes of (e) this synthetic excitable tissue in detail, including (e) dispersion relations, (e&eb) curvature-dependent wave front propagation, (e&eb) electrotonic coupling, and (e&eb) boundary effects.** |
| (253)**Authors then apply the biophysical characterization to (e) develop an optically reconfigurable bioelectric oscillator.** |
| (254)**These results demonstrate the feasibility of engineering bioelectric tissues capable of (e&eb) complex information processing with optical input and output. Received 12 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevX.6.031001 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society** |

**Scientists detect radio emission from a nearby brown dwarf July 6, 2016 by Tomasz Nowakowski report**

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| (255)**K2 and Spitzer images of W0607+24: The position of W0607+24 is marked by (e) a red circle. SDSS J060738.06+242939.5, marked by (e) an orange circle, are the brightest star in the K2 image.** |
| (256)**The three faint (r = 18.3 − 19.6) stars marked by cyan circles, contribute significant flux to (e) the K2 measurement but are (-) negligible at Spitzer's longer wavelength and (e &êb) higher resolution. Other stars with r < 20 in SDSS are marked with blue circles. Credit: Gizis et al., 2016.** |

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| (257)**(Phys.org)—Located some 23 light years away, the brown dwarf designated WISEP J060738.65+242953.4 or W0607+24 for short, turns out to be (=) a source of radio emission.** |
| (258)**According to a research paper published July 4, this substellar object showcases (eb) quiescent radio emission, making it one of (e) the most radio-faint, ultra-cool dwarfs yet detected.** |
| (259)**Brown dwarfs like W0607+24 are objects that are too large to be called planets and too small to be stars. With a mass below that necessary to maintain hydrogen-burning nuclear fusion reactions, they are much cooler and dimmer than (eb) main sequence stars.** |
| (260)**These objects could be also sources of (e) radio emission, but the cause of this process is still not completely understood.** |
| (261)**Due to its proximity, W0607+24 is (=) an excellent target for studying the nature of (e) brown dwarfs.** |
| (262) **It is the nearest known late-L dwarf in the northern hemisphere, and the third-nearest in (=) the whole sky, (eb) allowing researchers to study phenomena occurring in brown dwarfs.** |
| (263)**That's why a team of astronomers, led by John Gizis of the University of (e)Delaware, Newark, carried out a series of mid-infrared, radio, and optical observations of W0607+24 to study any activity in this curious substellar object.** |
| (264)**Motivated by the results from NASA's prolonged Kepler mission, named K2, which obtained (eb) long time-series photometry of W0607+24; they observed this object simultaneously with the Karl G. Jansky Very Large Array (VLA) in New Mexico and NASA's Spitzer Space Telescope.** |
| (265)**The most important results were provided by VLA, as the team found that W0607+24 is (=) a source of quiescent, non-bursting, radio emission.** |
| (266)**The scientists noted that the discovery of this process in this nearby object could be (=) crucial for our understanding of this process in (eb) ultracool dwarfs in general.** |
| (267)**"W0607+24 is a valuable test case for (e) understanding the mechanisms of radio emission in ultracool dwarfs.** |
| (268)**Some of the radio-detected objects emit (eb) rapid, intense bursts with high circular polarization, often found to occur periodically and attributed to coherent emission due to (e) the electron cyclotron maser instability," the paper reads.** |
| (269)**The astronomers also suggested that the relative radio faintness of W0607+24 may be due to (e) its orientation if (e) the emission is concentrated at the poles.** |
| (270)**However, follow-up observations are needed to better characterize (eb) the detected radio emission and to eliminate (e) all the uncertainties.** |
| (271)**The team also hopes for more significant findings when (e) studying this nearby substellar object in the future.** |
| (272)**"As the nearest known northern hemisphere late-L dwarf, this offers many possibilities for future studies. For example, it could serve as (=) an useful standard for (e) line-broadening and (e&eb) polarization studies.** |
| (273)**If there is a planetary system aligned with (e&eb) (the primary's rotation axis, then we would expect very little radial velocity signal but a strong astrometric signal," the researchers concluded.** |
| (274)**Moreover, the team plans to monitor (e&eb) additional late-M and L dwarfs in other K2 campaigns, expecting even more promising results.** |
| (275) **They are more distant than W0607+24, but because they are warmer, (eb) some are brighter and higher signal-to-noise than that described in this paper.** |

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| (276)**In addition to detecting the radio emission, (eb)the scientists also estimated the age, mass and radius of W0607+24** |
| (277)**. According to the study, the object is less than two billion years old, has (e) a maximum mass of about 0.055 solar masses and a radius approximately of 0.1 solar radii.** |
| (278)**More information: WISEP J060738.65+242953.4: A Nearby. Pole-On L8 Brown Dwarf with Radio Emission, arXiv: 1607.00943 [astro-ph.SR] arxiv.org/abs/1607.00943** |
| (279)**Authors present a simultaneous, multi-wavelength campaign targeting the nearby (7.2 pc) L8/L9 (optical/near-infrared) dwarf WISEP J060738.65+242953.4 in (eb) the mid-infrared, radio, and optical.** |
| (280)**Spitzer Space Telescope observations show no (e) variability at the 0.2% level over 10 hours each in the 3.6 and 4.5 micron bands.** |
| (281) **Kepler K2 monitoring over 36 days in Campaign 0 rules out (e) stable periodic signals in (eb) the optical with amplitudes great than 1.5% and periods between 1.5 hours and 2 days.** |
| (282)**Non-simultaneous Gemini optical spectroscopy detects lithium, constraining (e) this L dwarf to be less than ~2 Gyr old, but no Balmer emission is observed.** |
| (283)**The low measured projected rotation velocity (v sin i < 6 km/s) and lack of variability are (=) very unusual compared to other brown dwarfs, and authors argue that this substellar object is (=) likely viewed pole-on.** |
| (284)**Authors detect quiescent (non-bursting) radio emission with (e&eb) the VLA.** |
| (285) **Amongst radio detected L and T dwarfs, it has (e) the lowest observed L_nu and the lowest v sin i.** |
| (286)**They discuss the implications of pole-on detection for (e) various proposed radio emission scenarios. © 2016 Phys.org** |

**A new look at the galaxy-shaping power of black holes July 6, 2016**

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| (287)**It wasn't until the advent of X-ray astronomy that the full picture of (e) the galaxy-shaping power of black holes began to emerge (eb) with the ability to see plasma.** |
| (288)**Data from a now-defunct X-ray satellite is providing (eb) new insights into the complex tug-of-war between galaxies, the hot plasma that surrounds them, and the giant black holes that lurk in their centres.** |
| (289)**Launched from Japan on February 17, 2016, the Japanese space agency (JAXA) Hitomi X-ray Observatory functioned for just over a month before contact was lost and (e&eb) the craft disintegrated.** |
| (290)**But the data obtained during those few weeks was enough to paint (eb) a startling new picture of the dynamic forces at work within galaxies.** |
| (291)**New research, published in the journal Nature today, reveals data that shows just how important the giant black holes in galactic centres are to the evolution (eb) of the galaxies as a whole.** |
| (292)**"We think that supermassive black holes act(e&eb) like thermostats," said Brian McNamara, University Research Chair in Astrophysics at the University of Waterloo. "They regulate (e&eb) the growth of galaxies."** |

**Champagne bubbles of plasma**

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| (293)**During its brief life, the Hitomi satellite collected X-ray data from (e) the core of the Perseus cluster, an enormous gravitationally-bound grouping of (e) hundreds of galaxies.** |
| (294) **Located some 240 million light years from earth, the Perseus cluster is one of (e) the largest known structures in the universe.** |
| (295) **The cluster includes not only the ordinary matter that makes up the galaxies, but an (e) "atmosphere" of hot plasma with a temperature of tens of millions of degrees, as well as a halo of (e) invisible dark matter.** |
| (296)**Earlier studies, going back to the 1960s, have shown that each of the galaxies in the cluster - and indeed most galaxies - likely contains (e) a supermassive black hole in its centre, an object 100 million to more than ten billion times as massive as our sun** |

**New look at the galaxy-shaping power of black holes**

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| (297)**Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving.** |
| (298)**"These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science.** |
| (299) **"Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy."** |
| (300)**This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return.** |
| (301)**The remaining matter gets absorbed into the black hole, adding to (e) its mass.** |
| (302)**The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne.** |
| (303)**The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution.** |
| (304)**Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form.** |
| (305)**"Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara.** |
| (306)**"It creates (eb) these bubbles, and the bubbles keep the plasma hot.** |
| (307)**That's what prevents galaxies from (e) becoming even bigger than they are now."** |
| (308)**Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.** |
| (309)**In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.** |
| (310)**Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.** |

(311)**"It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.**

(312)**"It's like nature's thermostat, which keeps these galaxies from (e) growing.**

(313) **If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."**

**A new look at the galaxy-shaping power of black holes**

(314)**An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the**

(315)**McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.**

(316)**Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.**

**Doomed satellite missions**

(317)**Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.**

(318)**At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.**

(319)**Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.**

(320)**The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.**

(321)**These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.**

(322)**"But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people**

(323)**. It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."**

(324)**The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.**

(325) **"This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More**

**information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website**

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(326)**Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.**

(327)**They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.**

(328)**However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.**

(329) **Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.**

(330)**X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well.**

(331)**The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma.**

(332)**These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6.**

(333)**Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus.**

(334)**A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate.**

(335)**Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure.**

**Dark matter guides growth of supermassive black holes February 18, 2015**

(336)**This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy.**

(337)**Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole.**

(338) **But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy.**

(339)**A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole.**

(340) **It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth.**

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| (341)**"There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA).** |
| (342)**This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies.** |
| (343)**However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated.** |
| (344)**In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1.** |
| (345) **We know dark matter exists only from (e) its gravitational effects.** |
| (346) **It holds together galaxies and (e&eb) galaxy clusters.** |
| (347) **Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years.** |
| (348)**To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies.** |
| (349) **They used star motions as tracers to (e) weigh the galaxies' central black holes.** |
| (350) **X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto.** |
| (351)**They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone.** |
| (352)**This connection is likely to be related to (e&eb) how elliptical galaxies grow.** |
| (353) **An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together.** |
| (354)**Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole.** |
| (355)**"In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp** |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (356)**The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the …** |
| (357)**Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness?** |

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| (358)**Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness.** |
| (359)**"Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain."** |
| (360)**Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness.** |
| (361)**"With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication."** |
| (362)**Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories.** |
| (363)**The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen.** |
| (364) **In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer.** |
| (365)**Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants.** |
| (366)**These were treated as (=) "aware" and "unaware" trials, respectively.** |
| (367)**Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it.** |
| (368)**The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another.** |
| (369)**Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion.** |
| (370) **No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness.** |

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| (371)**"We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said.** |
| (372)**"Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness."** |
| (373)**The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner.** |
| (374)**Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain.** |
| (375)**"We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said.** |
| (376) **"This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients** |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

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| (377)**Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter.** |
| (378) **As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined.** |
| (379)**Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions.** |
| (380)**"Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.** |
| (381)**The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.** |
| (382)**However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.** |
| (383)**The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.** |
| (384) **Noda and his colleagues instead studied the change in emissions over (e&eb) time.** |

(385) **By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.**

(386)**Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).**

(387)**The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.**

(388)**Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.**

(389)**"The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp**

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(390)**The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking "palm tree" organization.**

(391)**The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society**

(392)**(Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible?**

(393) **With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem.**

(394)**This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other.**

(395)**According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible.**

(396)**This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges.**

(397) **In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century.**

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| (398)**Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching.** |
| (399)**In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom.** |
| (400) **Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth.** |

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| **Spheres on trees** |

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| (401)**Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration.** |
| (402)**They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement.** |
| (403) **Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum.** |
| (404)**By visualizing the problem in this way, the researchers noticed that these particular graphs don't (e) have lots of branches extending from lots of (e) other branches.** |
| (405) **Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure.** |
| (406)**The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network.** |
| (407)**As a result, it provides (eb) an important clue for finding the global minimum.** |
| (408)**It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum.** |
| (409)**This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large.** |

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| **Thompson's problem 2** |

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| (410)**By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society** |
| (411)**The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima.** |
| (412) **The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node.** |

**Implications for Thomson's and Smale's problems**

(413)**Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before.**

(414)**"Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around."**

(415)**As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem.**

(416)**"I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view."**

(417)**In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges.**

(418) **From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear.**

(419)**"While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK.**

(420)**"For larger numbers, there will likely be (=) multiple funnels.**

(421)**Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum."**

(422)**Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.**

(423)**"The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.**

(424)**"Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org**

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**Introduction:**

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price,**

**University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss.**

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| (1) | **Black Holes are regions of (e) space in which gravitational fields are so strong that no particle or signal can escape (eb) the pull of gravity.** |
| (2) | **The boundary of this no-escape region is (=) called the event horizon, since distant observers outside the black hole cannot (e) see (cannot get light from) events inside.** |
| (3) | **Although the fundamental possibility of such an object exists within (eb) Newton's classical theory of gravitation, Einstein's theory of gravity makes (eb)black holes inevitable under some circumstances.** |
| (4) | **Prior to the early 1960s, black holes seemed to be (=) only an interesting theoretical concept with no (e) astrophysical plausibility, but with (e&eb) the discovery of quasars in 1963 it became clear that very exotic astrophysical objects could exist.** |
| (5) | **Nowadays it is taken for granted that black holes do exist in (eb) at least two different forms.** |
| (6) | **Stellar mass black holes are (=) the endpoint of the death of some stars, and supermassive black holes are the result of (e) coalescences in the centers of most galaxies, including our own.** |
| (7) | **No signal can propagate from (e) inside a black hole, but the gravitational influence of a black hole is (=) always present.** |
| (8) | **This influence does not (e) propagate out of the hole; it is permanently present outside, and depends only on (=) the total amount of mass, angular momentum, and electric charge that have gone into forming (eb) the hole.** |
| (9) | **Black holes can be detected through the influence of this (e&eb) strong gravity on the surroundings just outside the hole.** |
| (10) | **In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them.** |
| (11) | **Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors.** |

## Classical vs. relativistic black holes

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| (12) | **Something like a black hole exists within (eb) Newton's classical theory of gravity.** |
| (13) | **In that theory, an energy argument tells us that there is (=) an escape velocity $v_{esc}=\sqrt{2GM/R}$ from (e) the surface of any spherical object of (e) mass M and radius R** |
| (14) | **If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity.** |
| (15) | **Thus the condition for such an "unseeable" object is $R<2GM/c2$** |
| (16) | **In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape.** |
| (17) | **This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity.** |
| (18) | **To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it.** |
| (19) | **The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c .** |

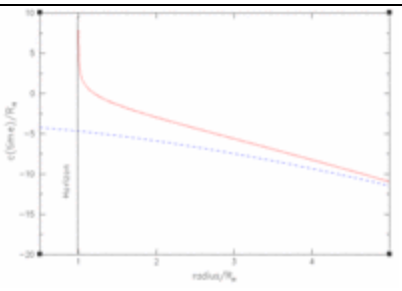
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| (20) This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed. |
| (21) An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity. |
| (22) Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable. |

**Stationary black holes**

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| (23) In the Newtonian theory, gravity is described by (e) the potential Φ. |
| (24) Inside a spherical object the form of (e) Φ(r) depends on (=) the interior structure, but in the vacuum outside (e) matter the potential −GM/r depends only on (=) the interior mass. |
| (25) Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object. |
| (26) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward. |
| (27) This horizon radius RH=2GM/c2 is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory. |
| (28) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution. |
| (29) Radius RHhere actually means that the area of the event horizon is (=) 4πR2H. |
| (30) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution. |
| (31) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity. |
| (32) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own. |
| (33) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole. |
| (34) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density. |
| (35) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed. |
| (36) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole. |
| (37) In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter. |
| (38) Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum. |
| (39) In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior. |

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| **(40) The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object.** |
| **(41) Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum.** |
| **(42) These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.]** |
| **(43) All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time.** |
| **(44) But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?"** |
| **(45) The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point.** |
| **(46) One of these coordinates is (=) called the "coordinate time."** |
| **(47) Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate.** |



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| **(48) Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole** |
| **(49) At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole.** |
| **(50) Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations.** |
| **(51) For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree** |
| **(52) An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events.** |
| **(53) The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving.** |
| **(54) Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle.** |

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| (55) The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return. |
| (56) But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon. |
| (57) Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future. |
| (58) The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?" |

**Black hole parameters**

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| (59) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin). |
| (60) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun. |
| (61) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits. |
| (62) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form. |
| (63) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds. |
| (64) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2. |
| (65) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed. |
| (66) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer. |
| (67) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2) |
| (68) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials. |
| (69) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable |
| (70) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite. |
| (71) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole. |
| (72) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars. |
| (73) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity. |
| (74) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity. |

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| **(75) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite.** |
| **(76) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity.** |
| **(77) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum.** |
| **(78) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship.** |
| **(79) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= (4πG2/c4)[(M+M2−a2−−−−−−−√)2+a2].** |
| **(80) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored.** |

**Dynamical black holes**

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| **(81) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape.** |
| **(82) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, a horizon can be born.** |
| **(83) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid.** |
| **(84) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward.** |
| **(85) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations.** |
| **(86) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes.** |
| **(87) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3.** |
| **(88) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/** |
| **(89) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers.** |
| **(90) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission.** |
| **(91) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.** |
| **(92) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.** |

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| **(93) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.** |
| **(94) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.** |
| **(95) The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature 6×10−8 (Mo/M) K.** |
| **(96) Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.** |

Then $[Y_6]_6[\gamma_6]_6, a_6) = E_6(a_6)e^{-\gamma_6 a_6}$. by a simple calculation,

$$\int_0^\infty \beta_6(a_6)Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma}\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}\ da\right);$$

$$\int_0^\infty Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma_6}\left(\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}\ da_6\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_6(0)}{\gamma_6}[f_{60}^6]\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}da_6\ (\int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6)$$

$$+\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right).$$

$$\left(1 + \frac{(\rho_0)_6(0)}{\gamma_6}(f_{60}^{'})_6\left(\int_0^\infty E_6(a_6)da_6 - \int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}da_6\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right)(1+\frac{(f_{60}^{'})_6}{\gamma_6}(\rho_0)_6(0)\int_0^{\infty}E_6(a_6)da_6)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_6(a_6)Y_6(\gamma_6 a_6)da_6=1-\int_0^{\infty}\beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\neq 0,$$

Since $\int_0^{\infty}\beta_6(a_6)Y_6(\gamma_6 a_6)\int_0^{a}\alpha_6(\tau_6)Y^{-6}(\gamma_6,\tau_6)d\tau_6 da_6>0$. Therefore, the characteristic solution is

$$\gamma_6=-(f_{60}^{'})_6(\rho_0)_6\int_0^{\infty}E_6(a_6)da_6.$$

Since $f_6'(P_0)_6>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_6),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE OF DEBIT TYPES AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_6),0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a),Q^*)$. It must satisfy the equations:

$[p_a^*]_6 + (\mu_6(a_6) + [f_6]_6(P_6^*) + [f_2]_6(Q_6^*))P_6^* = 0,$ (26)

$p_6^*(0) = h_6(Q_6^*)\int_0^\infty \beta_6(a_6)p_6^*(a_J)p_6^*(a_6)da_6$ (27)

$g_6(P_6^*) - b_6 = 0,$ (28)

Where $P_6^*=\int_0^\infty p_6^*(a_6)\,da_6$

Equation (28) yields $P^* = g^{-6}(b)$,

From (26),

$p_6^*(a) = p_6^*(0)e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6}$

Substituting into (27), we need to solve

$$\left(1 - h_6(Q_6^*)\int_0^\infty \beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6}da_6\right)p_6^*(0) = 0$$

For $Q_6^*$

Fix $p_6^*(0)$ and $P_6^*$ Let

$$H_6(z_6) = 1 - h_6(z_6)\int_0^\infty \beta_6(a_6)\,e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6}da_{,6}$$

And assume that

$$[k_3]_6 \int_0^{\infty} \beta_6(a_6) e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(g_6^{-6}(b_6))a_6 - [k_2]_6 a_6}\, da_6 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty} [H_6(z_6)] > 0.$$

Clearly,

$$\frac{dH_6}{dz_6} = h_6(z)([f_2^{'}]_6(z_6)) \int_0^{\infty} a_6 \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - f_6(P_6^*)a_6 - f_2(z_6)a_6} da_6$$

$-h^{'}(z_6) \int_0^{\infty} \beta_6(a_6)\, e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6}\ da_6$

By the assumptions on $\left[f_2^6\right]_6 and\ [h^6]^6 dH_6/dz > 0$. Hence, if

$$H_6(0) = 1 - \int_0^{\infty} \beta_6(a_6) e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6} da_6 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_6(0) < 0$ follows if $g_6^-(b_6) < [f_6^{-6}]_6(r_6)$, where $r_6$ is the solution of (12).

Indeed, since

$$R_6(s_6) = \int_0^{\infty} \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 e^{-y_6^a}} da_6$$

Is a decreasing function of $s_6$ when $P_6^* = [g^{-6}]_6(b_6) < [f_6^{-6}]_6(r_6)$, that is $[f_6]_6(P_6^*) < r$, $R\big(f_6(P^*)\big) > 1$which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_6(g^{-6}(b))) > 6, i.e.,$

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6\left(g_6^{-6}(b)\right)a_6} > 1. \qquad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_6^*(a_6), Q_6^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_6^*(a_6), Q_6^*)$ is more complicated. We again consider perturbations $m_6(t_6 a_6) = p_6(t_6, a_6) - p_6^*(a_6)$ and $W_6(t) = Q_6(t) - Q^*$.Then

$$M_6(t) \equiv \int_0^\infty m_6(t, a_6)da_6 = P_6(t) - P_6^*,$$

$$[m_6]_6 + [m_a]_6 = [p_6]_6 + [p_a]_6 - \left[p^*\right]_6(a_6) \approx -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6$$

$$-[f_{6*}^6]_6 p_6^*(a_6)M - [f_{2*}^6]_6 p_6^*(\mathrm{a})W_6,$$

$$m_6(t,0)p_6(t,0) - p_6^*(0) \approx h_* \int_0^\infty \beta_6(a_6)m_6(t,a_6)\,da + \frac{[h_*^6]_6}{[h_*]_6}\, p_6^*(0)W_6$$

$$\frac{dW_6}{dt} = \frac{dQ_6}{dt} \approx [g_*^6]_6 Q_6^* M_{,6}$$

Where $[f_{6*}]_6$ denoting $[f_6(P^*)]_6$, $[f_{6*}]_6$ denoting $[f_2(P^*)]_6$, $\left[f_{6*}^{'}\right]_6$ denoting $\left[f^{'}{}_6(P^*)\right]_6\left[f_{2*}^{'}\right]_6$ denoting $[f_2]_6$ $(Q_6^*,$ $h_{*}$ denoting $h_*^{'}$ denoting $h_6^{'}$ $(Q_6^*)$, and $\left[g_*^{'}\right]_6$

Denoting $\left[\dot{g}\right]_6(P_6^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_6]_6 + [m_a]_6 = -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6 - \left[\dot{f}_{6*}\right]_6[p_6^*](a_6)M_6 - \left[\dot{f}_{2*}\right]_6 p_6^*(a_6)W_6$$

$$m_6(t,0) = [h_*]_6 \int_0^\infty \beta_6(a_6)m_6(t,a_6)da_6 + \frac{[h_*^6]_6}{[h^*]_6}[p^*]_6(0)\, W_{,6} \qquad (30)$$

$$\frac{dW_6}{dt} = [g_*^6]_6[Q_6^*]M_6$$

Integrating $[m_6]_6(t,a_6)$ with respect to a from 0 to ∞, we have

$$\frac{dM_6(t)}{dt} = \int_0^\infty [mt]_6(t,a_6)da_6$$

$$= -\int_0^\infty [m_a]_6\,(t,a_6)da_6 - \int_0^\infty \mu(a_6)m(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + \left[\dot{f}_{6*}\right]_6 \int_0^\infty p_6^*(a_6)da_6\right)M_6 - [f_{2*}^6]_6 \int_0^\infty p_6^*(a_6)da_6 W_6$$

$$= m_6(t,0) - \int_0^\infty \mu_6(a_6)\, m_6(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + \left[\dot{f}_{6*}\right]_6 P_6^*\right)M_6 - \left[\dot{f}_{2*}\right]_6 P_6^* W_{\cdot 6}$$

From (30), we arrive at the equations

$$\frac{dW_6(t)}{dt} = \int_0^\infty ([h_*]_6\beta_6(a_6) - \mu_6(a_6))m_6(t,a_6)da_6$$

$$-\left([f_{6*}] + [f_{2*}]_6 + [f_{6*}]_6 P_6^*\right) M_6 + \left(\frac{[h_*']_6}{[h_*]^6}\, P_6^*(0) - f_{2*}' P_6^*\right) W_6,$$

(36)

$$\frac{dW_6(t)}{dt} [g'_*]_6\, Q_6^*\, M_{.6}$$

By the mean value theorem,

$$\int_0^\infty ([h_{"*}]_6 \beta_6(a_6) - \mu_6(a_6)) m_6(t, a_6) da_6 = ([h_*]_6 \beta_6(\theta_6) - \mu_6(\theta_6)) M_{,6}$$

Where $0< \theta_6 = \theta_6(t) < \infty$.

And if $a_6^2 [a_6^2]_6 < 4 [a_2]_6 < 4 a_2 e$,

$\text{Re}\,\lambda_6 [\lambda_6]_6 = \text{Re}\lambda_2 [\lambda_2]_6 = -\frac{6}{2} a_6 [a2\ ]_6 < 0.$

Let $\Delta = \max_{2<u<\infty} |h_* [h_*]_6 \beta_6(a_6) - \mu_6(a_6)|$. Then $|n_6(t)| \leq \Delta \forall t \geq 0$. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_6^2 < 4 a_2 e$ and $[a_6] > 2\Delta$,or $[a_6^2] \geq 4[a_2]e$ and $(a_6 - \sqrt{(a_6^2)_6 - 4(a_2)_6 e}) > 2\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+6}\right)$ and $\beta_6(a_6) = ae^{-ma}$ Then $\Delta = [\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_6(a_6)\mathrm{e}^{-\int_0^{\mathrm{a}} \mu_6(\tau_6) + d\tau_6 \mathrm{da}} = \frac{6}{(\mathrm{m}+\mu_0+1)^2} + \frac{2}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_6 + [\mu_0]$ 0.5. Assume$[f_6](x_6) = K_6 x_6$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_6(x_6) = [0x_6]/(1+x_6)$; the functional responses $[[f_2](x) = [0]\, x_6/(1+x_6)$and $h_6(x_6)$=(1+$\omega_6 x_6$)/$(1+x_6)$(See Holling [7] and Rashevsky [13])

Equation (64),now has the form of

$$\frac{m_6 + [\mu_0] + [\mathrm{K}][P_0])}{[m_6]_6 + (\mu_0)_6 + [\mathrm{K}][P_0])_6^3\ [\mathrm{K}]_6 [P_0]_6^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(6+P_0) < b_6$i e $[P_0] < b/(\delta_6 - b_6)$then the equilibrium$(P_0)$ $[P_0]_6(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_6$

$[f_{6*}] = K_6\ \mathrm{P}^*\ [f_{2*}] = \frac{\sigma \mathrm{Q}_6^*}{6+\mathrm{Q}_6^*} \quad [h_*] = \frac{6+\omega_6 \mathrm{Q}_6^*}{6+\mathrm{Q}_6^*}$

$[f_{6*}] = K_6\ [\mathrm{f}_{2*}^{6}] = \frac{\sigma_6}{(6+\mathrm{Q}_6^*)^2}[\ \mathrm{g}_*^{6}] = \frac{\sigma_6}{(6+\mathrm{P}_6^*)^2} \quad [\mathrm{h}_*^{6}] = \frac{\omega_6 - 6}{(6+\mathrm{Q}_6^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta$, and $\omega$ need to satisfy

$$\frac{\omega_6}{(m+[\mu_0]_6+1+\frac{K_6 b_6}{\delta_6+b_6}+\sigma)^2} + \frac{2\omega_6}{([m_6]_6+[\mu_0]_6+1+\frac{K_6 b_6}{\delta_6+b_6}+\sigma_6)^3} < 1$$

And

$$\frac{1}{(m_6 + [\mu_0]_6 + 1 + \frac{K_6 b_6}{\delta_6 - b_6})^2} + \frac{2\omega_6}{(m_6 + [\mu_0]_6 + 1 + \frac{K_6 b_6}{\delta_6 - b_6})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_6 + [\mu_0]_6 + 1 + \frac{K_6 b_6}{\delta_6 - b_6} + (m_6 + [\mu_0]_6 + \frac{K_6 b_6}{\delta_6 - b_6} + \sigma_6)\, Q_6^*)^3$$

$$= m_6 + [\mu_0]_6 + 3 + \frac{K_6 b_6}{\delta_6 - b_6}(1 + Q^*)(1 + \omega_6 Q_6^*$$

Now

$$\left[a_6^2\right] - 4\left[a_2\right] e = \left(\frac{2K_6 b_6}{\delta_6 - b_6}\right)^2 + \frac{4K_6 b_6 \sigma_6 Q_6^*}{(\delta_6 - b_6)(1 + Q_6^*)} + \frac{\sigma_6^2 Q_6^{*2}}{(1 + Q_6^*)^2}$$

$$\frac{4\, b_6 \sigma_6 (\delta_6 - b) Q^*}{\delta_6 (1 + Q_6^*)^2} - \frac{4\,(\delta_6 - b_6)^2 (1 - \omega_6) Q_6^*}{\delta_6 (1 + \omega_6 Q_6^*) 1 + Q_6^*)} P^*(0)$$

And

$$\Delta = \lim_{0 < a < \infty} |[h_*]_6\, ae_6^{mu}\ [\mu_0]_6 - \frac{a_6}{a_6 + 6}$$

Hence if $[a_6^2 > 4(a_2)_6 e$ and

$$[a_6] = \frac{2K_6 b_6}{\delta_6 - b_6} + \frac{\sigma_6 \, Q_6^*}{6 + Q_6^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_6]_{\,6} = [f_{6*}]_{\,6} + [f_{2*}]_{\,6} + [f_{6*}] \, P_6^*$$

$$[a_2]_{\,6} = \left[f_{2*}^{\,6}\right] P_6^* - \frac{h_*^6}{h_*} P^*(0)$$

$$e = [g_*^6]_{\,6} \, Q_6^*$$

$$\eta_6(t) = \left[h_*\right]_{\,6} \beta_6 \left(\theta_6(t)\right) - \mu_6(\theta(t)),$$

The system(36)can be written as the following system:

$$[M(t)]\left[a_6\right] M_6 - [a_2]_{\,6} W_6 - [n(t)] M_6$$

$$w_6(t) = e M_6$$

Write

$$A = \begin{pmatrix} -[a_6]_{\,6} - [a_2]_{\,6} \\ e \qquad\qquad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_6 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_6 \\ W_6 \end{pmatrix}$$

The solution of (32)is

$$X_6(t) = e^{A,t}[X]_6(0) + \int_0 e^{At}$$

Other problems to which the model is applied are: To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case $\beta\rightarrow\infty$ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness condition of (e) entropy. It is shown that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third

order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure authors have discovered. It is shown that Legendre invariant thermodynamic scalar curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws,

leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions, geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged AdS black holes by considering the case of the RN-AdS black holes surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase

transition. For the case ωq=−2/3 authors derive analytic expressions of (e) critical physical quantities, while (e) for cases ωq≠−2/3 appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory

with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=) chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations**

**Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions,

while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma? Paul M. Chesler STOP Paul M. Chesler, Krishna Rajagopal** calculate how the energy and (e&eb) the opening angle of jets in N=4 SYM theory evolve as they propagate through (e&eb) the strongly coupled plasma of that theory. **Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial

opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb) heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version) **Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb)

strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar (=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI:

10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills (SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as

well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as

will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations

in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh

reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the

quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of

chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstates thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11 August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and

represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there

are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.** Characterizing the behaviour of strongly coupled quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of

the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at λ=λc, the correlation length diverges as ξ∼|λ−λc|−ν, while (e) the gap between the ground state and first excited state vanishes as (e) ΔE∼ξ−z∼|λ−λc|zν. Exactly at the critical point λ=λc, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as ω∼|k−k0|z. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004 STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible

to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element. The basic architecture presented here can be expanded, offering (eb) the possibility for the coherent interaction of (e&eb) a large number of superconducting qubits. **Nature 449, 438-442 (27 September 2007) | doi:10.1038/nature06124; Received 18 April 2007; Accepted 25 July 2007 Coherent quantum state storage and transfer between two phase qubits via a resonant cavity Mika A. Sillanpää1, Jae I. Park1 & Raymond W. Simmonds1 STOP. Chui-Ping Yang, Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and**

**quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e) long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007**

# PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION ATTRITION MODELS

## MOODULE NUMBER SEVEN

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Systems to which have applied the model are as follows: Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures**

**have been proposed. An attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in (eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP. In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011**

**http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 Qubit transfer between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys. Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer. Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP. In recent years S. Lorenzo, T. J. G. Apollaro, S. Paganelli,**

**G. M. Palma, and F. Plastina P have investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys. Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP. Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2 [quant-ph] for this version) Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP. Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue from over (e&eb) a thousand plant and animal samples. NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP. Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this behavior, enhancing and accelerating (eb, eb+) the growth of correlations. Michael Kastner gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. Michael Kastner tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP. A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly**

**arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E. Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91, 075435 – Published 26 February 2015 STOP. Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1 proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction**

**involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point'' with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that (eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035 Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013 STOP Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and**

**cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen STOP Authors study the dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along**

**a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) Dynamic Service Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by H∗(X,(∧∗TX)⊗g), where X is the target hyperk\"ahler manifold. We further give a**

**mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition. Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb)**

**the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537 © 2016 American Physical Society Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant (Δα/α) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and Δα/α to (e&eb) this conclusion. However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts z≲3, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the Λ cold dark matter model value (jΛCDM=1); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in α(z), linked to (e&eb) a linear coupling (ζ) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order |ζ|∼10−4, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in Δα/α (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems,**

**authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP Authors measure the star formation properties of (e) two large samples of galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the star formation rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a function of environment. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]. Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size √ n × √ n? The problem is that √n time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time O (√n) for d ≥ 3, or O (√nlog5/2n) for d = 2. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ⁎This work was mostly done while the author was a PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship**

**and by CIAR. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e) sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP The relationship between the foundations of mathematics and (e&eb) physics is**

**a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb) flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph] (or arXiv: 1503.05600v1 [math-ph] for this version) Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time**

**dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to the fine structure constant. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor**

**goes to 0 as the (e&eb) time goes to 0, the big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph] (or arXiv: 1306.4613v1 [quant-ph] for this version) Effects of gauge theory based number scaling on geometry Paul Benioff STOP In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x, outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph] (or arXiv: 1211.3381v2 [quant-ph] for this version) Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value, \int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances,**

**the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include (e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP This work is based on (e) the field of reference frames based on (e) quantum representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology**

**(gr-qc) Journal reference: Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul Benioff STOP A quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference:**

**Journal of Physics: Conference Series 70 Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows: System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol**

## 5. MODEL FORMULATION

Let $p_7(t, a_7)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(g_7)Q_7$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_7(0) = 0, \quad g_7(P_7) > 0, \ P_7 \ 1, \qquad \lim_{P_7 \to \infty} g_7(P_7) = (K_7)_7 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_7 + (\rho_a)_7 + D_7(a_7, P_7, Q_7)\rho_7 + 0, \qquad 0 < a_7 \ \infty,$$

$$\rho_7(t, 0) = \int_0^\infty n_7(a_7, P_7, Q_7)\rho_7(t_7, a_7)da_7$$

Where $D_7$the mortality or debit type function is function and $n_7$ is the fertility function.

Put $_7$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_7}{dt}\right]_7 = x_7\left[\frac{dQ_7}{dt}\right]_7 - x_7\left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_7}{dt}\right]_7 = x_3\left[\frac{dQ_3}{dt}\right]_7 - x_4\left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_7}{dt}\right]_7 - x_6\left[\frac{dQ_3}{dt}\right]_7$$

Where $x_7, x_7, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_7$ and $g_7(P_7)Q_7$

Now when you add the term $(b_7Q_7)$ with $g_7(P_7)(Q_7)$ we are infect see is that $(b_7Q_7)$ is credited incrementally part of $g_7(P_7)(Q_7)$ or instant nearly, when subtracted from $\frac{dQ_7}{dt}$ , it means that $g_7(P_7)(Q_7)$ is slowly eaten by $\frac{dQ_7}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

Following equation throws light on important issues in the following cases

**Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA**

(425)(Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor.

(426)In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface.

(427)As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone.

(428) Such planets are all the "right" distance from (e) their star, and hopefully, also have water.

(429)But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone.

(430)For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement.

(431) But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway.

(432) But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain.

(433)It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life.

(434)Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface.

(435)The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113

**Abstract**

(436)The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question.

(437)Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry.

(438)This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry.

(439)Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments.

(440)Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy.

(441)The electronic and (e&eb) structural variability among them is extraordinary.

(442)The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations.

(443)The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface.

(444) The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers.

(445)Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp

**Water waves travel back in time to retrace their ripples By Jacob Aron (New Scientist)**

(446)Nudge your bath tub just right and (eb) you might end up reversing time.

(447)Don't worry – you won't end up wet, naked and without a towel in the Cretaceous period. Emmanuel Fort at the Langevin Institute in Paris, France, and his colleagues have developed what they call a "time mirror", which can make water waves rewind into (e&eb) their past. Humans and rubber duckies are (=) affected.

(448)To understand how this works, imagine dropping a stone in a pond and watching waves ripple out from a central point in an ever-expanding circle. The equations governing the behaviour of waves say that it should also be possible for a ring of ripples to refocus back into a single point, like a reversed movie of the dropped stone. But we never actually see this happening, because it is vanishingly unlikely that an outside-in ring would form spontaneously. Put it in reverse

(449)Fort and others have previously shown that it is possible to engineer the situation so the waves turn back on themselves by kitting out the edge of your pond with sensors and wave generators, effectively recording and then playing back the waves in reverse. The process has been done with water, sound and even microwaves, and is used to produce clearer ultrasound images for medical uses.

(450)But now the team has discovered that simply dropping the tub downwards is enough to send water waves scurrying in reverse, with no need for complex sensors. That's because the velocity of water waves depends on acceleration due to gravity. "If you do a sudden jolt, you change the effective gravity, and the velocity of the wave changes," says Fort.

(451)The disruption essentially resets the starting point of the wave. Once normal gravity is restored, it starts to travel both backwards and forwards. "Suddenly, you freeze the wave. When you release it again, it doesn't remember that it was going one way," says Fort.

(452)Ripples in the water returned exactly to their start point - the shape of the Eiffel Tower

(453)Ripples in the water returned exactly to their start point – the shape of the Eiffel Tower

**Déjà vu**

(454)To show that the backward-travelling waves converge exactly on their starting point, just as if you were rewinding time, the team conducted a series of experiments. In some, they simply created a circular ripple of waves by dipping a probe in the water, similar to dropping a stone in the pond. In others, they used

wave emitters in the shape of France, the Eiffel Tower and a smiley face to create complex wave-fronts. In each case, the reversed waves converged to recreate their original shape.

(455)The experiments aren't truly reversing time, because the waves continue to travel out again once they have converged at the central point, points out Fort. "A real time-reversal would have a sink that absorbs all the waves," he says.

(456)Now the team hopes to conduct similar experiments with other kinds of waves, like light or sound, although it may prove more difficult to find an analogue of their time-mirror jolt in other wave mediums, he says.

(457)"The work is spectacular," says Geoffrey Edelmann of the US Naval Research Laboratory in Washington, D.C. "They essentially change the rules that govern how the wave propagates and, in doing so, create an entirely new way to focus waves free of spatial constraints." He also thinks the technique could be extended to other kinds of waves, and perhaps have medical applications such as lithotripsy, which uses ultrasound waves to break up kidney stones. Journal reference: Nature Physics, DOI: 10.1038/nphys3810

**Time reversal and holography with spacetime transformations Vincent Bacot, Matthieu Labousse, Antonin Eddi, Mathias Fink & Emmanuel Fort Nature Physics (2016) doi: 10.1038/nphys3810**

(458)Wave control is usually performed by (e) spatially engineering the properties of a medium.

(459) Because time and space play similar roles (e&eb) in wave propagation, manipulating time boundaries provides (eb) a complementary approach.

(460) Here, authors experimentally demonstrate the relevance of this concept by (e) introducing instantaneous time mirrors.

(461)Authors show with water waves that a sudden change of (e) the effective gravity generates (eb) time-reversed waves that refocus at the source.

(462)They generalize this concept for (e) all kinds of waves, introducing (eb) a universal framework which explains (eb) the effect of any time disruption on (e&eb) wave propagation.

(463)Authors show that sudden changes of the medium properties generate (eb) instant wave sources that emerge (eb) instantaneously from the entire space at (eb) the time disruption.

(464)The time-reversed waves originate from (e) these 'Cauchy sources', which are (+) the counterpart of Huygens virtual sources on (eb) a time boundary. It allows us to revisit the holographic method and introduce a new approach for wave control.

**Viewpoint: A Bird's Eye View of Circuit Photons Sebastian Schmidt, Institute for Theoretical Physics, ETH Zurich, 8093 Zurich, Switzerland June 30, 2016• Physics 9, 74 A scanning probe detects the quantum states of photons in a microwave circuit, providing the information needed for quantum simulations. Figure caption expand figure A. Houck/Princeton (deleted)**

(465)Figure 1: Scanning defect microscopy provides a map of photons in a resonator lattice. Houck and colleagues demonstrated the technique using (e) 49 resonators (grey lines) that were coupled together to (e&eb) form a kagome lattice.

(466)Quantum mechanics rules the dynamics of (e) light and matter.

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| (467) Yet performing a quantum-mechanical simulation of a material from first principles is (=) practically impossible on a classical computer because (e) the complexity of the simulation increases (eb) exponentially with the number of particles involved. |
| (468) The solution, according to Richard Feynman, was to build a machine out of (e) quantum building blocks that could directly emulate (e&eb) the material itself [1]. |
| (469)Prototypes of such quantum simulators that are based on (e) ultracold atoms, ions, photons, and superconducting microwave circuits are now available [2], with (e&eb) the latter, in particular, having attracted (e) Silicon Valley's interest. |
| (470)The challenge with these circuit-based simulators, however, is that they are (=) 2D, which complicates the readout of their constituent elements. Andrew Houck from Princeton University, New Jersey, and colleagues have now delivered an attractive solution by developing a technique [3], called scanning defect microscopy (Fig. 1), that determines (eb) the number of photons occupying (e) each mode of a 2D microwave circuit. |
| (471)It is this information that would serve as (=) the fundamental input and output for (e) certain quantum simulations. |
| (472)Superconducting microwave circuits combine electronic and (e&eb) photonic degrees of freedom [4, 5]. |
| (473)The main element of the circuit is (=) a transmission line, which is made up of (e) a central superconducting wire separated by a gap from (e) two grounded plates. |
| (474)All of these structures are on a single plane, as if one had taken a 2D slice through (e&eb) a coaxial cable. |
| (475)When truncated, the transmission line becomes (=) a resonator, which can host (e&eb) discrete photon modes within (eb) its gaps |
| (476) Large lattices of resonators can be engineered in various 1D or 2D geometries by coupling (e&eb) two, three, or more resonators together via (e&eb) a capacitive interface. In many ways, photons in such devices behave similarly to electrons in a solid. |
| (477)To make microwave circuits that can simulate (e&eb) quantum phenomena faster than a classical computer, however, resonator lattices have to be integrated with (e&eb) superconducting qubits |
| (478). Such qubits are controlled with (e&eb) electrical currents in a Josephson tunnel junction, and in many respects, they behave (e&eb) like artificial atoms, which couple to (e&eb) the photons in the circuit. |
| (479)As a result, superconducting microwave circuits can be used to (e) explore the coupling between the quantum states of light and (e&eb) matter, the regime of (e) circuit quantum electrodynamics (cQED). |
| (480)Photons in these devices often exhibit (eb) striking matter-like behavior [6, 7], providing (eb) the basis for the simulation of complex materials. |
| (481) Such circuits can be fabricated on a substrate using (e) standard lithographic techniques, with (e&eb) qubits and resonators that are (=) hundreds of micrometers or even millimeters in size. |
| (482)Small-scale realizations of cQED lattices have already been used for (e) proof-of-principle quantum simulations [8, 9]. |
| (483)However, future large-scale simulations will have to be initialized and measured (eb) efficiently—steps known as (=) drive and readout. |
| (484) This requires connecting the lattices to (e) external circuitry, such as voltage sources, circulators, and amplifiers, which typically reside outside (e) the chip. |
| (485)So far, this connectivity has been realized through (e&eb) external transmission lines, but these elements can only couple to (e&eb) the resonators at the edges of (e) the lattice because of (e) its 2D nature. |

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| (486)Such a side view provides (eb) only limited information about (e&eb) the quantum state and spatial correlations of photons inside (eb) the bulk resonators, seriously limiting (e) the applicability of circuit QED devices for (e) quantum simulations. |
| (487)The probe developed by Houck and colleagues, scanning defect microscopy (SDM), solves (e) this problem by (e) moving outside the plane of the chip to provide (eb) a fuller image of photons within (eb) the chip. |
| (488) In SDM, a small sapphire tip is positioned a few tens of (e) micrometers above the surface of one of the resonators (Fig. 1). |
| (489)The presence of this dielectric modifies (e&eb) the evanescent electric field from the circuit photons, which extends (eb) just outside the plane of the chip. |
| (490) Specifically, the tip increases (eb) the local capacitance of the transmission line, downshifting (e) the local resonator frequency. |
| (491) This in turn causes (eb) a drop in the transmission of photons through (e&eb) the lattice when (e) it is driven at its nominal resonance frequency. |
| (492)Houck's team showed that if the sapphire tip has only a weak (linear) effect on the resonator frequency, then (eb) a simple analytic relationship exists between the induced frequency shift, (e&eb) the change in transmission, and (e&eb) the local photonic occupation of each resonator. |
| (493)This relationship provides (eb) a recipe for detecting the local photon occupancy inside (eb) the lattice. |
| (494)First, calibrate the probe by determining (eb) the frequency shift for (e) a single isolated resonator as (=) a function of the tip's position. |
| (495) Second, measure the transmission through (e&eb) the whole lattice with and without (e&eb) the probe. |
| (496)Finally, use the measured transmission drop with (e&eb) the probe present to find (eb) the photonic occupation of (e) an individual resonator in (eb) the lattice. |
| (497)A map of the photon-mode occupancies can be obtained by (e) scanning the probe across the whole lattice. |
| (498)Houck's team used (e) SDM to image the occupancies of (e) the photon modes of (e) 49 coupled resonators arranged in (eb) a special geometry called a kagome lattice (see details in Fig. 1). |
| (499) (To simplify the measurement, no superconducting qubits were (=) incorporated with (e&eb) the resonators.) |
| (500)The occupancies they measured were in excellent agreement with (=) theoretical predictions, corresponding to (e&eb) a fidelity of 0.99. |
| (501) The high fidelity is even more impressive considering the weak signal-to-noise ratio of (e) their technique (which is limited by (e) the intentionally small size of the frequency-shifted region) and (e&eb) the technical challenge of (e) positioning the probe inside (eb) a cryostat for (e) in situ measurements. |
| (502) While SDM shares similarities with (e&eb) so-called bead-pull experiments for accelerators and scanning gate microscopy for (e) nanostructures, it is truly a (=) new tool for quantum electronics. |
| (503) First application of SDM is likely to be (=) the simulation of 2D topological phenomena and (e&eb) geometrically frustrated lattices (of which the kagome lattice is (=) one example.) |
| (504)Both are associated with (e&eb) macroscopically degenerate systems, which are (=) notoriously difficult to simulate on (e&eb) a classical computer. |
| (505)Such circuits are (=) genuinely open systems (photons can be added or lost). |
| (506)This combined with (e&eb) their inherent strong light-matter coupling, makes (eb) such quantum simulators ideal candidates for (e) the study of many-body, out-of-equilibrium phenomena, such as (=) thermalization or nonequilibrium phase transitions. |

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| (507) SDM could also lead to (eb) the discovery of new exotic phenomena, which are (=) unknown in traditional solid-state systems. |
| (508)In a nutshell, SDM constitutes a new and indispensible tool for (e) future experiments on quantum simulations with microwave photons. This research is published in Physical Review X. |

**Imaging Photon Lattice States by Scanning Defect Microscopy D. L. Underwood, W. E. Shanks, Andy C. Y. Li, Lamia Ateshian, Jens Koch, and A. A. Houck Phys. Rev. X 6, 021044 – Published 30 June 2016**

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|---|
| (509)Microwave photons inside (eb) lattices of coupled resonators and superconducting qubits can exhibit (eb) surprising matter like behavior. |
| (510) Realizing such open-system quantum simulators presents (eb) an experimental challenge and requires (e) new tools and measurement techniques. |
| (511) Here, authors introduce scanning defect microscopy as one such tool and illustrate its use in mapping (e&eb) the normal-mode structure of microwave photons inside (eb) a 49-site kagome lattice of coplanar waveguide resonators |
| (512) Scanning is accomplished by (e) moving a probe equipped with a sapphire tip across (e&eb) the lattice. |
| (513)This locally perturbs (e&eb) resonator frequencies and induces (eb) shifts of the lattice resonance frequencies, which authors determine by (e) measuring the transmission spectrum. |
| (514) From the magnitude of mode shifts, author can reconstruct (eb) photon field amplitudes at (eb) each lattice site and thus create (eb) spatial images of the photon-lattice normal modes. Received 4 December 2015 DOI:http://dx.doi.org/10.1103/PhysRevX.6.021044 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society |

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|---|
| |
| We assume that |
| $$D_7 = \mu_7(a_7) + (f_7)_7(P_7)_7 + (f_7)_7(Q)_7 \quad (7)$$ |
| $$n_7 = \beta_7(a_7)h_7(Q_7), \quad (3)$$ |
| Where $\mu_7$the age- dependent debit rate whit is $\mu_7(a_7)$ ⍰ $\forall a > 0$; $(f_7)_7$ is the density- dependent death rate with $f_7(0) = 0$, $(f_7^7)_7(P_7)0$, $\forall P$ ⍰0, and $\lim_{p\to\infty}(f_7)_7(P_7) = (k_7)_7$ ⍰∞; $f_7$ is the functional response of creditors to prey with $[f_7]_7(0) = 0, (f_7^7)_7(Q) > 0, \forall Q$ ⍰0, and $\lim_{Q\to\infty}(f_7)_7(Q_7) = [k_7]$ ⍰∞; $\beta$ is the age dependent birth rate with $\beta(a)$ ⍰$0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^7(Q)$ ⍰$0, \forall Q$ ⍰0, and $\lim_{Q\to\infty} h(Q) = k_3$ ⍰0. |

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_7 + (\rho_a)_7 = -\mu_7(a_7) + (f_7)_7 P_7) + \big(f_7)_7(Q_7)\big)(\rho_7)_7 \quad (4)$$

$$\rho_7(t_7, 0) = h_7(Q_7) \int_0^\infty \beta_7(a_7)\rho_7(t, a_7)da, \quad (5)$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(P_7)Q_7 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_7$ and $Q_7$ of (4-6) satisfy

$$0 \leq \rho_7(t, a_7) \leq \rho_7^*(t, a_7),$$

$$0 \leq Q_7(t) \leq Q_7^*(t),$$

Where $\rho_7^*(t, a_7)$ and $Q_7^*(t)$, are solutions of

$$(\rho_t^*)_7 + (\rho_a^*)_7 = -(\mu_7(a_7) + (f_7)_7(P_7^*) + (f_7)_7 Q_7^*))\,\rho_7^*, \quad (7)$$

$$\rho_7^*(t, 0) = \int_0^\infty \beta_7(a_7)\,\rho_7^*(t, a_7)da_7 \quad (8)$$

$$\frac{dQ_7^*}{dt} = -b_7 Q_7^* + g_7(P_7^*)Q_7^*, \quad (9)$$

With $P_7^*(t) = \int_0^\infty \rho_7^*(t, a_7)da_7$

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From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_7^*(t)}{dt} = P_7^*(t)\big(r_7 - (f_7)(P_7^*) - (f_7)_7(Q^*)\big), \quad (10)$$

$$\frac{dQ_7^*(t)}{dt} = Q_7^*(t)(-b_7 + g_7(P_7^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_7(a)e^{-\int_0^a \mu_7(\tau)dt}e^{-ra}\,da_7 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_7^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(t)dt}\,da_7 > 1, \quad (H1)$$

Which guarantees $r_7 > 0$.

*Remark.* If $\beta_7(a) = [\beta_0]_7 a_7 e^{-\alpha_7 a_7}, \mu_7(a_7) \equiv \mu_7$ then the solution of (12) is

$$r_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_7]_7(P_7) = 0, [f_7](Q_7) = \lambda_7 Q_7\ h_7(Q_7) = 1, \quad and\ \ g_7(P_7) = c_7 P_7$$

$$[\rho_7]_7 + [\rho_a]_7 = -(\mu_7 + \lambda_7 Q_7)\rho_7$$

$$\rho_7(t,0) = \int_0^{\infty} [\beta_0]_7 a_7 e^{-\alpha_7 a_7} \rho_7(t,a_7)\, da_7$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + c_7 P_7 Q_7$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_7}{dt} = \omega_7 P_7 - \lambda_7 P_7 Q_7$$

$$\frac{dQ_7}{dt} = b_7 Q_7 - c_7 P_7 Q_7$$

Where $\omega_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$

## 6. THE EQUILIBRIUM $(\rho_0)_7(a_7), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_7(a_7), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_7 + (\rho_a)_7 = -\big(\mu_7(a_7) + (f_7)_7(P_7)\big)\rho,$$

$$\rho_7(t,0)$$

$$= \int_0^\infty \beta_7(a_7)\rho_7(t,a_7)da_7 \qquad (13)$$

The solution of (13) has the form of

$$\rho_7(t,a_7) = Ce_7^{r(t-a_7)} - \int_0^a (\tau)d\tau - \int_0^7 (f_7)_7\big(P_7(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_7(t) = C_7 e^{r_7 t - \int_0^t (f_7)_7(P_7(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_7^{-7}(r)$ if $k_7 > r$. This is equivalent to

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{(k_7)_7 a_7} da_7 < 1.$$

In fact, if we define

$$R_7(s_7) = \int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau)d\tau} e^{-s_7 a_7}\, da_7$$

$R(s)$ is a decreasing function. $R_7(r_7) = 1$ And $R_7(k_7)_7 < 1$ imply $r_7 < (k_7)_7$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_7(a_7), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{k_7 a_7} da_7$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_7 \equiv \int_0^\infty (\rho_0)_{77}(a_7)da_7$ *satisfies*

$$\int_0^{\infty} \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{-(f_7)_7(P_0)a_7} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_7(a_7), 0)$

Let $u_7(t, a_7) = \rho(t, a_7) - \rho_0(a_7)$ Then

$$N_7(t) \equiv \int_0^{\infty} u_7(t, a_7) da_7 = \int_0^{\infty} \rho_7(t, a_7) da_7 - \int_0^{\infty} (\rho_0)_7(a_7) da = P_7(t) - P_0)_7$$

And

$$(u_7)_7 + (u_a)_7 = (\rho_t)_7 + (\rho_a)_7 - (\rho_0^7)_7(a_7)$$

$$= -\mu_7(a_7) + (f_7)_7(P_7) + (f_7)_7(Q_7))(u_7 + (\rho_0)_7) + \mu_7(a_7)(\rho_0)_7 + (f_7)_7(P_0)_7(\rho_0)_7$$

$$= -\mu_7(a_7)u_7 - (f_7)_7(P_7) + (f_7)_7(Q_7))\, u_7 - (\rho_0)_7\, (f_7)_7(P_7) - (f_7)_7(P_0)_7$$

$$+ (f_7)_7(Q_7)), \qquad (15)$$

Where $(\rho_0^7)_7(a_7)$ denotes $d(\rho_0)_7(a_7)/da_7$

Expansion $(f_7)_7(P_7)$ and $(f_7)_7(Q_7)$ at $(P_0)_7$ and 0 yield respectively

$$(f_7)_7(P_7) = (f_7)_7((P_0)_7 + N_7)$$

$$= (f_7)_7(P_0)_7 + (f_7^7)_7(P_0)_7 N_7 + o_7(N_7), \qquad (16)$$

$$(f_7)_7(Q_7)$$

$$= (f_7^7)_7(0)Q + o_7(Q_7) \qquad (17)$$

Write $f_7(P_0) = f_{70}, f_7'(P_0) = f_{70}'$, and $f_7'(0) = f_{70}'$ and substitute (16) and (17) into (15) to obtain

$$(u_7)_7 + (u_a)_7 \approx -(\mu_7(a_7) + (f_{70})_7)u - (f_{70}^7)_7(\rho_0)_7(N_7) - (f_{70}^7)_7(\rho_0)_7(Q_7)$$

$$u_7(t, 0) = \rho_7(t, 0)(\rho_0)_7(0)$$

$$= h_7(Q_7)\int_0^{\infty} \beta_7(a_7)(u_7(t,a) + (\rho_0)_7(a_7)da_7 - (\rho_0)_7\,(0)$$

$$= h_7(Q_7)\int_0^{\infty} \beta_7(a_7)(u_7(t,a_7)da_7 + (h_7(Q_7) - 1)(\rho_0)_7\,(0) \qquad (18)$$

Expanding

$$h_7(Q_7) = 1 + \ (h^7)_7(0)Q_7 + 0(Q_7) \equiv 1 + h_0^{'}Q_7 + 0Q_7$$

And substituting it into (18), we arrive at

$$u_7(t,0) \approx \int_0^{\infty} \beta_7\,(a_7)u_7(t,a_7)da_7 + (h_{a_7}^{'})(\rho_0)_7Q_7$$





The linearization is the following :

$$[u_7] + [u_7] = -(\mu_7(a_7) + [t_{70}]_7u_7 - [\alpha_7]_7(a_7)N_7 - [\alpha_7]_7(a_7)Q_7$$

$$u_7(t,0) = \int_0^{\infty} \beta_7(a_7)u_7(t,a_7)da_7 + [\xi_7]Q_7 \qquad (19)$$

$$\frac{dQ_7}{dt} = -b_7Q_7$$

where

$$[\alpha_7]_7(a_7) = [t_7^7a]_7[p,o]_7(a_7)d2(a_7) = [f_{70}^7]_7(\rho\ o\ )_7(a_7)\xi = [\mathrm{h\ o}^7]_7[\rho\mathrm{o}]_7(\mathrm{o}) \text{ and } \mathrm{b}_7 = \mathrm{b}_7 - g_7(P_7)$$

If $g_7(P_0) > b$, clearly, $(\rho_0)_7(a_7), 0)$ is unstable and the predator population is persistent. Assume $g_7(P_0)_7 < b$.

Then, $Q_7(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_7 + u_a = -(\mu_7(a_7) + (f_{70})_7)u_7 - (\alpha_7)_7(a_7)N_7$$

$$= -(\mu_7(a_7) + f_{70})u_7 - \alpha_7(a_7)\int_0^\infty u_7(t, a_7)da_7$$

$$u_7(t, 0) = \int_0^\infty \beta_7\,(a_7)u_7(t, a_7)da_7$$

Using the method developed by Cushing [2, 3], we suppose

$$u_7(t, a_7) = \tilde{u}_7(a_7)^{e^{\gamma(t-a_7)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_7}\,da_7$. Then,

A solution of the first differential equation for $\tilde{u}(a_7)$ gives,

$$\tilde{u}_7(a_7) = e^{-\int_0^a \mu_7(\tau_7)d\tau - (f_{70})_7 a}\left(\tilde{u}_7(0) - w\int_0^a e^{\int_0^\tau \mu_7(s_7)ds + (f_{70})}\tau_7\,(\alpha_7)_7(\tau_7)e^{\gamma_7\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_7(a_7)e^{-\gamma_7 a_7 - (f_{70})_7 a_7 - \int_0^a \mu_7(\tau_7)d\tau}\right)\tilde{u}_7(0)$$

$$+w\int_0^\infty \beta_7(a_7)e^{-\gamma_7 a_7 - (f_{70})_7 a_7 - \int_0^a \mu_7(\tau_7)d\tau}\int_0^a (\alpha_7)_7(\tau_7)e^{-\gamma_7\tau_7}e^{(f_{70})_7\tau_7 + \int_0^\tau \mu_7(s_7)ds}d\tau_7\,da_7$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty} e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau} da\right)(\tilde{u}_7(0))_7$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau}\int_0^a e^{(f_{70})_7\tau_7+\int_0^{\tau}\mu_7(s_7)ds+y\tau}\alpha_7(\tau)d\tau\, da\right)w=0. \quad (23)$$

Define $Y(\gamma,a)=e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(t_7)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_7\,(a_7)Y_7(\gamma_7,a_7)da_7\right)u_7(0)$$

$$+w\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\,da_7=0, \quad (24)$$

$$\left(\int_0^{\infty}Y_7(\gamma_7,a_7)da_7\right)\tilde{u}_7(0)$$

$$+\left(-1-\int_0^{\infty}Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\,da_7\right)w=0.$$

Equation (20) has a solution of the form of $u_7(t,a_7)=u_7(a_7)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_7(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_7(\gamma_7,a_7)=\int_0^a(\rho_0)_7(\tau_7)Y^{-7}\,(\gamma_7,a_7)d\tau_7$. We obtain the characteristic equation

$$f_{70}^{l}{}_{7}\left(\int_0^{\infty} Y_7(\gamma_7,a_7)da_7\right)\left(\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$+\left(1-\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)da\right)\left(1+(f_{70}^{l})_7\int_0^{\infty}Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_7(a_7)=(\rho_0)_7(0)e^{-\int_0^a \mu_7(\tau_7)d\tau_7-(f_{70})_7a_7}$$

$$I_7(\gamma_7,a_7)=\int_0^a(\rho_0)_7(\tau_7)e^{\gamma_7\tau_7+(f_{70})_7\tau+\int_0^{\tau}\mu_7(s_7)ds_7}d\tau$$

$$=\frac{(\rho_0)_7(0)}{\gamma_7}(e^{\gamma_7a_7}-1)$$

Let

$$E_7(a_7)=e^{-(f_{70})_7a_7-\int_0^a\mu_7(\tau_7)d\tau_7}$$

Then $[Y_7]_7[\gamma_7]_7,a_7)=E_7(a_7)e^{-\gamma_7a_7}$. by a simple calculation,

$$\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da=\frac{[\rho_0]_7(0)}{\gamma}\left(1-\int_0^{\infty}\beta_7(a_7)E_7(a_7)e^{-\gamma_7a_7}\ da\right);$$

$$\int_0^{\infty}Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da=\frac{[\rho_0]_7(0)}{\gamma_7}\left(\int_0^{\infty}E_7(a_7)e^{-\gamma_7a_7}\ da_7\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_7(0)}{\gamma_7}[f_{70}^7]\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\,(\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7)$$

$$+\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right).$$

$$\left(1+\frac{(\rho_0)_7(0)}{\gamma_7}(f_{70}^{'})_7\left(\int_0^\infty E_7(a_7)da_7-\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right)(1+\frac{(f_{70}^{'})_7}{\gamma_7}(\rho_0)_7(0)\int_0^\infty E_7(a_7)da_7)=0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)da_7 = 1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7 \neq 0,$$

Since $\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)\int_0^a \alpha_7(\tau_7)Y^{-7}(\gamma_7,\tau_7)d\tau_7 da_7 > 0$. Therefore, the characteristic solution is

$$\gamma_7 = -(f_{70}^{'})_7(\rho_0)_7\int_0^\infty E_7(a_7)da_7.$$

Since $f_7'(P_0)_7 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_7), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE OF DEBIT AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_7), 0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_7 + (\mu_7(a_7) + [f_7]_7(P_7^*) + [f_7]_7(Q_7^*))P_7^* = 0, \quad (76)$$

$$p_7^*(0) = h_7(Q_7^*)\int_0^\infty \beta_7(a_7)p_7^*(a_J)p_7^*(a_7)da_7 \quad (77)$$

$$g_7(P_7^*) - b_7 = 0, \quad (78)$$

Where $P_7^* = \int_0^\infty p_7^*(a_7)\, da_7$

Equation (28) yields $P^* = g^{-7}(b)$,

From (26),

$$p_7^*(a) = p_7^*(0)e^{\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(Q_7^*)a_7}$$

Substituting into (27), we need to solve

$$\left(1-h_7(Q_7^*)\int_0^{\infty}\beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7-[f_7]_7(P_7^*)a_7-[f_7]_7(Q_7^*)a_7}da_7\right)p_7^*(0)=0$$

For $Q_7^*$

Fix $p_7^*(0)$ and $P_7^*$ Let

$$H_7(z_7)=1-h_7(z_7)\int_0^{\infty}\beta_7(a_7)\,e^{-\int_0^a \mu_7(\tau_7)d\tau_7-[f_7]_7(P_7^*)a_7-[f_7]_7(z_7)a_7}da_{,7}$$

And assume that

$$[k_3]_7\int_0^{\infty}\beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7-[f_7]_7(g_7^{-7}(b_7))a_7-[k_7]_7a_7}\,da_7<1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_7(z_7)]>0.$$

Clearly,

$$\frac{dH_7}{dz_7}=h_7(z)(\,[f_7']_7(z_7)\,)\int_0^{\infty}a_7\beta_7(a_7)\,e^{-\int_0^a \mu_7(\tau_7)d\tau_7-f_7(P_7^*)a_7-f_7(z_7)a_7}da_{7}$$

$-h'(z_7)\int_0^{\infty}\beta_7(a_7)\,e^{\int_0^a \mu_7(\tau_7)d\tau_7-[f_7]_7(P_7^*)a_7-[f_7]_7(z_7)a_7}\quad da_7$

(See author's It is a debit-credit world)

By the assumptions on $\left[f_7^7\right]_7 and\ [h^7]^7 dH_7/dz > 0$. Hence, if

$$H_7(0) = 1 - \int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7} da_7 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_7(0) < 0$ follows if $g_7^-(b_7) < [f_7^{-7}]_7(r_7)$, where $r_7$ is the solution of (12).

Indeed, since

$$R_7(s_7) = \int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 e^{-y\frac{a}{7}}} da_7$$

Is a decreasing function of $s_7$ when $P_7^* = [g^{-7}]_7(b_7) < [f_7^{-7}]_7(r_7)$, that is $[f_7]_7(P_7^*) < r$, $R\left(f_7(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

**THEOREM 3**

Assume that the hypothesis (H3) holds and $R(f_7(g^{-7}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7\left(g_7^{-7}(b)\right)a_7} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_7^*(a_7), Q_7^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_7^*(a_7), Q_7^*)$ is more complicated. We again consider perturbations $m_7(t_7 a_7) = p_7(t_7, a_7) - p_7^*(a_7)$ and $W_7(t) = Q_7(t) - Q^*$.Then

$M_7(t) \equiv \int_0^\infty m_7(t,a_7)da_7 = P_7(t) - P_7^*,$

$$[m_7]_7 + [m_a]_7 = [p_7]_7 + [p_a]_7 - [p^*]_7(a_7) \approx -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7$$

$$-[f_{7*}^7]_7 p_7^*(a_7)M - [f_{7*}^7]_7 p_7^*(\mathrm{a})W_7,$$

$$m_7(t,0)p_7(t,0) - p_7^*(0) \approx h_* \int_0^\infty \beta_7(a_7)m_7(t,a_7)\,da + \frac{[h_*^7]_7}{[h_*]_7}\, p_7^*(0)W_7$$

$$\frac{dW_7}{dt} = \frac{dQ_7}{dt} \approx [g_*^7]_7 Q_7^* M,_7$$

Where $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}']_7$ denoting $[f'_7(P^*)]_7[f_{7*}']_7$ denoting

$[f_7]_7$ $(Q_7^*,$ $h_{7}^*$ denoting $h_*'$ denoting $h_7'$ $(Q_7^*)$, and $[g_*']_7$

Denoting $[g']_7(P_7^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$[m_7]_7 + [m_a]_7 = -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7 - [f'_{7*}]_7[p_7^*](a_7)M_7 - [f_{7*}']_7 p_7^*(a_7)W_7$

$m_7(t,0) = [h_*]_7 \int_0^\infty \beta_7(a_7)m_7(t,a_7)da_7 + \frac{[h_*^7]_7}{[h^*]_7}[p^*]_7(0)\,W,_7$ (30)

$$\frac{dW_7}{dt} = [g_*^7]_7[Q_7^*]M_7$$

Integrating $[m_7]_7(t,a_7)$ with respect to a from 0 to ∞, we have

$$\frac{dM_7(t)}{dt} = \int_0^\infty [mt]_7(t, a_7)da_7$$

$$= -\int_0^\infty [m_a]_7\,(t, a_7)da_7 - \int_0^\infty \mu(a_7)m(t, a_7)da_7$$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f'_{7*}]_7 \int_0^\infty p_7^*(a_7)da_7\right)M_7 - [f_{7*}^7]_7 \int_0^\infty p_7^*(a_7)da_7 W_7$$

$= m_7(t, 0) - \int_0^\infty \mu_7(a_7)\, m_7(t, a_7)da_7$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 - [f_{7*}^{'}]_7 P_7^* W_{.7}$$

From (30), we arrive at the equations

$$\frac{dW_7(t)}{dt} = \int_0^\infty ([h_*]_7\beta_7(a_7) - \mu_7(a_7))m_7(t, a_7)da_7$$

$$-\left([f_{7*}] + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 + \left(\frac{[h_*']_7}{[h_*]^7}\, P_7^*(0) - f_{7*}' P_7^*\right)W_7,$$

(31)

$$\frac{dW_7(t)}{dt}[g'_*]_7\, Q_7^*\, M_{.7}$$

By the mean value theorem,

$$\int_0^\infty ([h_{"*}]_7\beta_7(a_7) - \mu_7(a_7))m_7(t, a_7)da_7 = ([h_*]_7\beta_7(\theta_7) - \mu_7(\theta_7))M,_7$$

Where $0< \theta_7 = \theta_7(t) < \infty$.

And if a $_7^7$ [a $_7^7$]2<4 [a $_7$]2<4 a $_7$e,

Re $\lambda_7[\lambda_7]_7 = \mathrm{Re}\lambda_7[\lambda_7]_7 = -\frac{7}{7}a_7[\mathrm{a}2\ ]_7 < 0$.

Let ⊿=max$_{7<u<\infty}$ |h $_*$ [h $_*$] $_7$β $_7(a_7) - \mu_7(a_7)$|. Then|$n_7$(t)| ⊿∀t . Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium(ρ $^*$ (a), $\boldsymbol{Q}^*$) is asymptotically stable if either a $_7^7$ <4 a $_7$ e and [a $_7$] > 1⊿,or [a $_7^7$] [a $_7$]e and (a $_7 - \sqrt{(\mathrm{a}_7^7)_7 - 4(\mathrm{a}_7)_7\mathrm{e}} > 1⊿$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{\mathrm{a}}{\mathrm{a}+7}\right)$ and β $_7(a_7)$ = ae-$^{\mathrm{ma}}$ Then ⊿=[$\mu_0$] + 2 Simple calculation leads to

$$\int_0^\infty \beta_7(a_7)\mathrm{e}^{-\int_0^{\mathrm{a}}\mu_7(\tau_7)+d\tau_7\mathrm{da}} = \frac{1}{(\mathrm{m}+\mu_0+1)^7} + \frac{1}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_7$ + [$\mu_0$] .5. Assume[$f_7$]($x_7$)= $K_7x_7$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_7(x_7)$= [0$x_7$]/(2 + $x_7$); the functional responses [[$f_7$]($x$) = [0] $x_7$/(1 + $x_7$)and $h_7(x_7)$=(1+$\omega_7x_7$)/(1 + $x_7$)(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_7 + [\mu_0] + [\mathrm{K}][P_0])}{[m_7]_7 + (\mu_0)_7 + [\mathrm{K}][P_0])_7^3\ [\mathrm{K}]_7[P_0]_7^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1+P_0) < b_7$ i e $[P_0] < b/(\delta_7 - b_7)$then the equilibrium$(P_0)$ $[P_0]_7(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_7$

$[f_{7*}] = K_7\ \mathrm{P}^*\ [f_{7*}] = \frac{\sigma \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*} \quad [h_*] = \frac{7+\omega_7 \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*}$

$[f_{7*}] = K_7\ [\mathrm{f}_{7*}^{\ 7}] = \frac{\sigma_7}{(7+\mathrm{Q}_7^*)^7} [\ \mathrm{g}_*^{\ 7}] = \frac{\sigma_7}{(7+\mathrm{P}_7^*)^7} \quad [\mathrm{h}_*^{\ 7}] = \frac{\omega_7 - 7}{(7+\mathrm{Q}_7^*)^7}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_7}{(m + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma)^7} + \frac{1\omega_7}{([m_7]_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma_7)^3} < 1$$

And

$$\frac{1}{(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^7} + \frac{1\omega_7}{(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $\mathrm{Q}^*$ which can be solved from the following equation:

$(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7} + (m_7 + [\mu_0]_7 + \frac{K_7 b_7}{\delta_7 - b_7} + \sigma_7)\ \mathrm{Q}_7^*)^3$

$= m_7 + [\mu_0]_7 + 3 + \frac{K_7 b_7}{\delta_7 - b_7}(1 + \mathrm{Q}^*)(1 + \omega_7 \mathrm{Q}_7^*$

Now

$$\left[\mathrm{a}\begin{matrix}1\\1\end{matrix}\right]-4\left[\mathrm{a}_{7}\right]\mathrm{e}=\left(\frac{1K_7 b_7}{\delta_7-b_7}\right)^{7}+\frac{4K_7 b_7\sigma_7\,\mathrm{Q}_{7}^{*}}{(\delta_7-b_7)(7+\mathrm{Q}_{7}^{*})}+\frac{\sigma_{7}^{7}\mathrm{Q}_{7}^{*7}}{\left(1+\mathrm{Q}_{7}^{*}\right)^{7}}$$

$$\frac{4\,b_7\sigma_7\,(\delta_{7}-\mathrm{b})\mathrm{Q}^{*}}{\delta_{7}\,(1+\mathrm{Q}_{7}^{*})^{7}}-\frac{4\,(\delta_{7}-b_7)^{7}\,(1-\omega_7)\mathrm{Q}_{7}^{*}}{\delta_{7}\left(1+\omega_7\mathrm{Q}_{7}^{*}\right)(1+\mathrm{Q}_{7}^{*})}\mathrm{P}*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_{7}\,ae_{7}^{mu}\;\;[\mu_0]_{7}-\frac{a_7}{a_7+7}$$

Hence if $[\mathrm{a}_{7}^{7}>4(a_7)_{7}\mathrm{e}$ and

In the following are given some more modular structures that are applied with the model and the concomitant results derived thereof:

Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments

(1) We have an intuitive understanding of how the world works. Our intuition works well in our everyday lives, but it doesn't quite match (e&eb) reality.

(2) One of the basic assumptions authors make about objects is that they are (=) objectively real.

(3) For example, when you leave your apartment in the morning, you imagine your house is right where you left it as you go about your day. Barring some calamity, you expect to find your apartment just as you left it when (e) you return in the evening.

(4) This idea is known as Macrorealism, in that a macroscopic object such as your house is still there even when (e) you're not looking at it

(5) In scientific terms we would say the state of an object exists independent of (e) an observer. While this seems obviously true, we can test the idea experimentally.

(6) To show this, let's consider an experiment. Suppose we have a box that contains a coin. At any time we can open the box and (eb) see whether it's heads or tails.

(7) The catch is (=) that the box has to follow a certain rule. Whatever you observe initially, (eb) you must observe the same thing ten seconds later.

(8) If you open the box and see heads, then (eb) ten seconds later when you open the box you must also see heads.

$$[a_7] = \frac{1K_7 b_7}{\delta_7 - b_7} + \frac{\sigma_7 \, \mathrm{Q}_7^*}{1 + \mathrm{Q}_7^*} > 1\vartriangle$$

This equation is important in its consequence in respect of following cases for which the model is applied:

**Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments**

(1) For the experiment we'll open the box, then open it five seconds later, then finally open it again at ten seconds. We know the first and last observations have to agree, but what about the observation in the middle? It depends upon what the box does with (e&eb) the coin.

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| (2) One possibility is that the box does nothing, so if you see tails the first time (eb), you will see tails both other times. |
| (3) Another possibility is that each time you open the box (eb) it flips the coin when you close it. |
| (4) So the first time you see heads, (eb) the second time tails, and (eb) the third time heads again. |
| (5) But they could also make a tricky box that flips (e&eb) the coin at seemingly random times, but does so in such a way that every ten seconds it repeats (eb) the pattern. In that case, authors would always see the same thing the first and last time, but the middle observation would seem (eb) random. |
| (6) If Macrorealism is true, then (eb) middle observations being always the same as (=) the first and last observation, always opposite or always random are (=) the only possibilities for the experiment |

**Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry Download as PDF (Size: 62KB) HTML PP. 50-53 DOI: 10.4236/jqis.2011.12007 5,508 Downloads 14,712 Views Citations Author(s) Leave a comment Mohamed S. El Naschie**

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| (1) Building upon the pioneering work of J. Bell [1] and an incredible result due to L. Hardy [2] it was shown (eb) that the probability of quantum entanglement of two particles is (=) a maximum of 9.0169945 percent [2]. |
| (2) This happens to be exactly the golden mean to (e) the power of five (?5) [3-7]. |
| (3) Although it has gone largely unnoticed for a long time, this result was essentially established independently in a much wider context by (e) the present author almost two decades ago [3-6]. |
| (4) The present work gives two fundamentally different derivations of (e) Hardy's beautiful result leading to (eb) precisely the same general conclusion, namely that by virtue of (e) the zero measure of the underlying |
| (5) Cantorian-fractal spacetime geometry the notion of (e) spatial separability in quantum physics is (=) devoid of any meaning [7]. The first derivation is purely logical and uses (e) a probability theory which combines the discrete with (e&eb) the continuum. |
| (6) The second derivation is purely geometrical and topological using (e) the fundamental equations of a theory developed by the author and his collaborators frequently referred to as E-infinity or Cantorian spacetime theory [3-7]. KEYWORDS Hardy's Quantum Entanglement, Golden mean, Cantor sets, Fractal Spacetime, E-Infinity Theory, Quantum Mechanics, J. S. Bell Cite this paper M. El Naschie, "Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry," Journal of Quantum Information Science, Vol. 1 No. 2, 2011, pp. 50-53. Doi: 10.4236/jqis.2011.12007. |

**Chaos, Solitons & Fractals Volume 28, Issue 4, May 2006, Pages 857–878 Sir Hermann Bondi 1919–2005 Cantorian spacetime and Hilbert space: Part I—Foundations G. Iovane http://dx.doi.org/10.1016/j.chaos.2005.08.074**

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| (1) Authors are going to show the link between the ε(∞) Cantorian space and( e&eb) the Hilbert spaces H(∞). |
| (2) In particular, El Naschie's ε (∞) is (=) a physical spacetime, i.e. an infinite dimensional fractal space, where (e) time is spacialized and the transfinite nature manifests (eb) itself. |
| (3) El Naschie's Cantorian spacetime is an arena where (e) the physics laws appear at each scale in a self-similar way linked to (e&eb) the resolution of the act of observation. |
| (4) By contrast the Hilbert space H (∞) is (=) a mathematical support, which describes (eb) the interaction between the observer and (e&eb) the dynamical system under (e&eb) measurement. |

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| (5) The present formulation, which is based on (e) the non-classical Cantorian geometry and topology of spacetime, automatically solves (e) the paradoxical outcome of (e) the two-slit experiment and the so-called particle–wave duality. |
| (6) In particular, measurement (i.e. the observation) is equivalent to (=) a projection of ε(∞) in the Hilbert space built on 3 + 1 Euclidean spacetime. |
| (7) Consequently, the wave–particle duality becomes (=) a mere natural consequence of (e) conducting an experiment in (eb) a spacetime with (e&eb) non-classical topological and (e&eb) geometrical structures, while (e) observing and taking measurements in (eb) a classical smooth 3 + 1 Euclidean spacetime. |
| (8) In other words, the experimental fact that a wave–particle duality exists is (=) an indirect confirmation of the existence of ε(∞) and a property of (e) the quantum–classical interface. |
| (9) Another direct consequence of the fact that real spacetime is) =) infinite dimensional hierarchical ε(∞) is (=) the existence of scaling law R(N), introduced by the author, which generalizes the Compton wavelength. |
| (10)It gives an answer to the problem of (e) segregation of matter at different scales, and shows (eb) the role of fundamental constants such as (=) the speed of light and Plank's constant h in the fundamental lengths scale without (e) invoking the principles of quantum mechanics. Copyright © 2005 Elsevier Ltd. All rights reserved |

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**Chaos, Solitons & Fractals Volume 14, Issue 9, December 2002, Pages 1341–1351 Fractal strings as an alternative justification for El Naschie's Cantorian spacetime and the fine structure constant Carlos Castro, http://dx.doi.org/10.1016/S0960-0779(02)00110-8**

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| (1) Beginning with the most general fractal strings/sprays construction recently expounded in the book by Lapidus and Frankenhuysen, it is shown how the complexified extension of El Naschie's cantorian-fractal spacetime model belongs to (e&eb) a very special class of families of fractal strings/sprays whose scaling ratios are given by (e) suitable pinary (pinary, p prime) powers of the Golden Mean. |
| (2) They then proceed to show why the logarithmic periodicity laws in (eb) Nature are direct physical consequences of (e) the complex dimensions associated with (e&eb) these fractal strings/sprays. |
| (3) Authors proceed with a discussion on quasi-crystals with (e&eb) p-adic internal symmetries, von Neumann's Continuous Geometry, the role of wild topology in (e&eb) fractal strings/sprays, (e&eb) the Banach-Tarski paradox, and tesselations of (e) the hyperbolic plane, quark confinement and (e&eb) the Mersenne-prime hierarchy of (e) bit-string physics in determining the fundamental physical constants in Nature. Fax: +1-914-965-3397 Copyright © 2002 Published by Elsevier Ltd. About ScienceDirectRemote access Shopping cart Contact and support Terms and conditions Privacy policy Cookies are used by this site. For more information, visit the cookies page. Copyright © 2016 Elsevier B.V. or its licensors or contributors. Science Direct ® is a registered trademark of Elsevier B.V. |

**Chaos, Solitons & Fractals Volume 21, Issue 2, July 2004, Pages 261–271 On Cantorian spacetime over number systems with division by zero Jakub Czajko, http://dx.doi.org/10.1016/j.chaos.2003.12.046**

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| (4) Division by zero is proposed in order to open (eb) new horizons for (e) abstract mathematics and physics. |
| (5) Although the formerly forbidden operation may be of (e) little practical value, the feasibility of unrestricted operations is needed for (e) meaningful deployment of linear vector spaces, whose duals are (=) inverse transposes of their primary vector spaces. |
| (6) Both quantum mechanics and (e&eb) pan-geometry of multispatial hyperspace operate on (e&eb) vectors that belong to (e&eb) mutually dual linear vector spaces. |
| (7) The division by zero also may remove (e) many years old paradox related to (e&eb) multiplication of zeros and makes it possible to introduce (eb) the concept of bigroup as (=) an additive and multiplicative group at the same time. |
| (8) This operational extension to (e&eb) algebraic groups makes thus (eb) an infinite-dimensional, fractal Cantorian spacetime (=) the preferred, invertible, operationally unrestricted, abstract mathematical infrastructure for physics. Tel.: +1-925-681-2058 Copyright © 2003 Elsevier Science Ltd. All rights reserved. |

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The nontrivial equilibrium is stable.

Some of the examples in which stability analysis is of seminal importance are: These are the system to which the

model becomes applicable

**Observation of a Dissipation-Induced Classical to Quantum Transition J. Raftery, D. Sadri, S. Schmidt, H. E. Türeci, and A. A. Houck Phys. Rev. X 4, 031043 – Published 8 September 2014**

(111)Here, authors report the experimental observation of (e) a dynamical quantum phase transition in (eb) a strongly interacting open photonic system.

(112) The system studied, comprising a Jaynes-Cummings dimer realized on (eb) a superconducting circuit platform, exhibits (eb) a dissipation-driven localization transition.

(113)Signatures of the transition in the homodyne signal and photon number reveal (eb) this transition to be from (e) a regime of classical oscillations into (e&eb) a macroscopically self-trapped state manifesting revivals, (=) a fundamentally quantum phenomenon.

(114) This experiment also demonstrates (eb) a small-scale realization of a new class of quantum simulator, whose well-controlled (e&eb) coherent and dissipative dynamics is suited to (e) the study of quantum many-body phenomena out of (e) equilibrium. Received 10 December 2013 DOI:http://dx.doi.org/10.1103/PhysRevX.4.031043 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**Viewpoint: Fermionic Vortices Find their Dual John McGreevy, Department of Physics, University of California at San Diego, La Jolla, CA 92093-0319, USA June 27, 2016• Physics 9, 72Theoretical work reveals a surprising relationship between the physics of fermionic vortices and quantum electrodynamics. Figure caption expand figure APS/Alan Stonebraker Figure 1: The action of discrete symmetries on the fermionic particle (top) and its dual fermionic vortex (bottom).**

(115)Dirac fermions in 2+1 dimensions (two spatial dimensions plus time) have been found in (eb) graphene and on (e&eb) the surfaces of topological insulators (TIs).

(116) A trio of papers now shows that these particles can also be found in (eb) a dual description of quantum electrodynamics (QED) [1–3].

(117)This deceptively simple statement joins (e&eb) together several disparate lines of progress in theoretical physics: interacting topological insulators, (e&eb) anomalies in quantum field theory, (e&eb) spin liquids, and (e&eb) compressible Quantum Hall States.

(118)Humans are not very good at understanding interacting many-body systems.

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| (119)We understand the behavior of free particles and can sometimes use (e) perturbation theory to (e) incorporate weak interactions between (e&eb) them. |
| (120)But what happens when we couple (e&eb) the particles strongly? |
| (121)To reveal the interesting collective phenomena that can result, we've had to rely largely on (e) experiments. |
| (122)Examples include (e) superconductivity (e&eb) superfluidity (e&eb) the integer and (e&eb) fractional quantum Hall effects (e&eb) the fractionalization of charge and spin, and (e&eb) the wetness of water. |
| (123)The fractional quantum Hall effect (FQHE) is a physical phenomenon in which the Hall conductance of 2D electrons shows (eb) precisely quantised plateaus at (eb) fractional values of (e) {\displaystyle e^{2}/h} e^{2}/h. |
| (124) It is a property of (e) a collective state in which electrons bind (e&eb) magnetic flux lines to make (eb) new quasiparticles, and excitations have (e) a fractional elementary charge and possibly also fractional statistics. The 1998 Nobel Prize in Physics was awarded to Robert Laughlin, Horst Störmer, and Daniel Tsui for the discovery and explanation of the fractional Hall effect.[1] |
| (125) However, Laughlin's explanation was (=) a phenomenological guess and only applies to (e&eb) fillings {\displaystyle \nu =1/m} \nu =1/m where (e) {\displaystyle m} m is an odd integer. The microscopic origin of the FQHE remains unknown and is a major research topic in condensed matter physics. (Wikipedia) |
| (126)The quantum Hall effect (or integer quantum Hall effect) is a quantum-mechanical version of (e) the Hall effect, observed in (eb) two-dimensional electron systems subjected to (e&eb) low temperatures and strong magnetic fields, in which the Hall conductance σ undergoes (e&eb) certain quantum Hall transitions to take on (e&eb) the quantized values {\displaystyle \sigma ={\frac {I_{\text{channel}}}{V_{\text{Hall}}}}=\nu \;{\frac {e^{2}}{h}},} \sigma = \frac{I_\text{channel}}{V_\text{Hall}} = \nu \; \frac{e^2}{h}, where {\displaystyle I_{\text{channel}}} I_\text{channel} is the channel current, {\displaystyle V_{\text{Hall}}} V_\text{Hall} is the Hall voltage, e is the elementary charge and h is Planck's constant. The prefactor ν is known as the "filling factor", and can take on either integer (ν = 1, 2, 3, ...) or fractional (ν = 1/3, 2/5, 3/7, 2/3, 3/5, 1/5, 2/9, 3/13, 5/2, 12/5, ...) values. |
| (127) The quantum Hall Effect is referred to as (=) the integer or fractional quantum Hall effect depending on whether ν is an integer or fraction, respectively. |
| (128)The striking feature of the integer quantum Hall Effect is (=) the persistence of the quantization (i.e. the Hall plateau) as (e) the electron density is varied. |
| (129) Since the latter remains constant when (e) the Fermi level is in a clean spectral gap, it must be that the Fermi level is immersed in (eb) essential spectrum (i.e. of finite density of states), though Anderson localized. |
| (130) The quantization and stability of the Hall conductance in such conditions was achieved in 1994 by (e) Bellissard, van Elst and Schulz-Baldes using the methods of (e) Non-Commutative Geometry.[1] |
| (131)The fractional quantum Hall effect is more complicated, as its existence relies fundamentally on (e) electron–electron interactions. |
| (132)Although the microscopic origins of the fractional quantum Hall effect are unknown, there are several phenomenological approaches that provide (eb) accurate approximations. |
| (133) For example the effect can be thought of as (=) an integer quantum Hall effect, not of electrons but of (e) charge-flux composites known as (=) composite fermions. |

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| (134) In 1988, it was proposed that there was quantum Hall effect without Landau levels. This quantum Hall Effect is referred to as (=) the quantum anomalous Hall (QAH) effect. |
| (135)There is also a new concept of the quantum spin Hall effect which is an analogue of the quantum Hall effect, where (e) spin currents flow instead of (e&eb) charge currents.[2] |
| (136)In two dimensions, when classical electrons are subjected to (e&eb) a magnetic field they follow (eb) circular cyclotron orbits. |
| (137)When the system is treated quantum mechanically (eb) these orbits are quantized. |
| (138)The energy levels of these quantized orbitals take on (e) discrete values |
| (139):{\displaystyle \textstyle E_{n}=\hbar \omega _{c}(n+1/2),} \textstyle E_{n}=\hbar \omega _{c}(n+1/2), where ωc = eB/m is the cyclotron frequency |
| (140) These orbitals are known as Landau levels, and at weak magnetic fields, their existence gives rise to (eb) many interesting "quantum oscillations" such as the Shubnikov–de Haas oscillations and the de Haas–van Alphen effect (which is often used to map the Fermi surface of metals). |
| (141) For strong magnetic fields, each Landau level is (=) highly degenerate (i.e. there are many single particle states which have the same energy En). |
| (142)Specifically, for a sample of area A, in magnetic field B, the degeneracy of each Landau level is (=) {\displaystyle \textstyle N=g_{s} BA/\phi _ {0},} \textstyle N=g_{s} BA/\phi _ {0}, where gs represents (eb) a factor of 2 for spin degeneracy, and ϕ0 = 2·10−15 Wb is the magnetic flux quantum. |
| (143) For sufficiently strong B-fields, each Landau level may have (e) so many states that all of the free electrons in the system sit in only a few Landau levels; it is in this regime where (e) one observes the quantum Hall Effect. |
| (144)The integers that appear in the Hall Effect are (=) examples of **topological quantum numbers**. |
| (145) They are known in mathematics as the first **Chern numbers** and are closely related to (e&eb) **Berry's phase.** |
| (146)A striking model of much interest in this context is the Azbel-Harper-Hofstadter model whose quantum phase diagram is (=) the Hofstadter butterfly shown in the figure. (deleted due to spatial constraints) |
| (147)The vertical axis is (=) the strength of the magnetic field and the horizontal axis is the chemical potential, which fixes the electron density. The colors represent the integer Hall conductance. Warm colors represent positive integers and cold colors negative integers. The phase diagram is fractal and has structure on all scales. In the figure there is an obvious self-similarity. |
| (148)Concerning physical mechanisms, impurities and/or particular states (e.g., edge currents) are important for (e) both the 'integer' and 'fractional' effects. |
| (149) In addition, Coulomb interaction is also essential in (eb) the fractional quantum Hall effect. |
| (150)The observed strong similarity between integer and (=) fractional quantum Hall effects is explained by (e) the tendency of electrons to form bound states with (e&eb) an even number of magnetic flux quanta, called composite fermions. |
| (151)Duality is the circumstance in which cranking up interactions results in (eb) another system we understand, usually one that can be described in terms of (e&eb) another set of weakly coupled particles. |
| (152)When such a dual representation is found, (eb) it gives a solution to the strong coupling problem. |
| (153) A further reason that theoretical physicists get excited about dualities is that they undermine (e) the very notion of an “elementary particle.” |
| (154)In terms of the original variables, the dual, weakly coupled particles are (=) generally solitons—large, heavy, classical objects made up of (e) many quanta. |

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| (155)The remarkable thing is that solitons can become (=) "elementary," in the sense that the model that describes the system can be reformulated with the solitons as (=) the basic constituents. |
| (156)An example of a duality, which we will use below, arises in (eb) Maxwell's theory of electromagnetism, where (e) the strength of the interactions is determined by (e) the electric charge e. |
| (157)The duality exchanges electric and (e&eb) magnetic point charges, and the interaction strength in the dual description is 1⁄e, which is small when (e) e is large. |
| (158)This idea has been generalized to (e) many other field theories, thanks, in particular, to tools of (e) supersymmetry [4]. |
| (159)Another classic example of duality involves (e&eb) vortices. |
| (160) A vortex is a localized object in the plane around which a phase field winds (e&eb). (A phase field takes values on the circle.) |
| (161)Dualities relating particles and vortices have (e) a rich history in particle physics and condensed-matter physics. |
| (162) The simplest version relates the particles in the normal state of a superfluid to (e&eb) the vortices in a dual superconductor [5]. |
| (163)This charge-vortex duality has been (=) a workhorse in the study of two-dimensional condensed-matter systems, such as (=) fractional Quantum Hall States, superconducting thin films, and beyond-Landau critical phenomena [6]. |
| (164)Vortices can be (=) fermions (as they frequently are, for instance, in type-II superconductors). |
| (165) A dual description of a system with (e&eb) fermionic vortices thus requires (e) a new ingredient, since (e) it must include (e) such a fermionic particle. |
| (166)The new work [1–3] builds on (e&eb) earlier theoretical work [7] to make (eb) an appealing proposal for (e) the required duality: the fermionic vortex of QED in (eb) 2+1 dimensions is (=) a free Dirac fermion. The argument for this proposal builds on our understanding of TIs. |
| (167)TIs are distinguished only by (e) their edge states. |
| (168) The most famous example is a (3+1)-dimensional time-reversal invariant TI, which can host boundary states whose band structure is defined by (e) a single Dirac cone. |
| (169)A single time-reversal-symmetric Dirac cone is not possible without (e) the bulk TI, so its appearance in experiments provides (eb) a signature of the bulk phase. |
| (170) But the Dirac cone is not (e) the only possible signature. |
| (171) Deformations of the boundary conditions can dramatically change (e&eb) the spectrum of the edge states, yet preserve some essential "TI-ness." This essential property is (=) called an anomaly in high-energy theory. |
| (172)Max Metlitski of the University of California, Santa Barbara, and Ashvin Vishwanath of the University of California, Berkeley [1], and independently, Chong Wang and Todadri Senthil at the Massachusetts Institute of Technology, Cambridge [2], use this ambiguity in the edge theory and the electric-magnetic duality of an auxiliary theory of bulk electromagnetism coupled to (e&eb) the TI to argue for the duality. |
| (173)When I said that (3+1)-dimensional QED had a duality that interchanged the electric and magnetic charges, one could have (e) complained that the electron is (=) a fermion, while (e) a magnetic monopole may not be. |
| (174)However, the authors observe that in a TI, a charge-two monopole is (=) also a fermion. |

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| (175) By analyzing the spectrum of electric and magnetic charges, they conclude that the duality exchanges the electron and (e&eb) the charge-two monopole and, moreover, exchanges a TI with (e&eb) time-reversal symmetry with (e&eb) a "chiral TI," which instead has (e) particle-hole symmetry |
| (176). But this is just what is required of (e) a symmetry that exchanges a fermionic particle with (e&eb) a fermionic vortex: time-reversal symmetry maps the electron to itself but maps the vortex to (e&eb) another vortex whose phase winds (e&eb) the other way (see Fig. 1). |
| (177)The exchange of particles and holes, in turn, preserves (eb) the vortex. |
| (178)The bulk duality implies (eb) a duality between edge theories. |
| (179)The freedom to deform the edge of the TI implies (eb) a weak form of equivalence: although the edge theories are (=) topologically the same (in terms of symmetry assignments and (e&eb) anomalies), they need not have (e) identical low-energy physics. |
| (180)One side of the duality is (=) a free, massless Dirac particle, but the long-distance physics of (2+1)-dimensional QED is (=) a long-standing problem, which is important in the study of spin liquids. |
| (181) Many gauge theories in two dimensions develop (eb) a gap in their energy spectrum because of (e) the proliferation of monopoles. |
| (182) In the realization proposed by Metlitski and Vishwanath [1] and by Wang and Senthil [2], this cannot happen because (e) the monopoles carry electric charge, and indeed the stronger notion of duality as (=) an equivalence of low-energy physics may hold. |
| (183)Evidence for the stronger relation comes from (e) David Mross and colleagues at the California Institute of Technology, Pasadena [3], who have derived a lattice version of the duality using (e) a clever but explicit change of variables. |
| (184)Their construction is particularly interesting because (e) it uses (e) a lattice of one-dimensional models which themselves can only occur at the edge of a TI. |
| (185)Authors explicitly derive the duality between a free electronic Dirac cone and (e&eb) quantum electrodynamics in (eb) (2+1) dimensions (QED3) with (e&eb) N=1 fermion flavors. |
| (186)The duality proceeds via (e&eb) an exact, nonlocal mapping from electrons to (e&eb) dual fermions with (e&eb) long-range interactions encoded by (e) an emergent gauge field. |
| (187)This mapping allows (eb) us to construct (eb) parent Hamiltonians for exotic topological-insulator surface phases, derive (eb) the particle-hole-symmetric field theory of (e) a half-filled Landau level, and nontrivially constrain (e) QED3 scaling dimensions. |
| (188)They similarly establish duality between bosonic topological insulator surfaces and (e&eb) N=2 QED3. |

**Black hole thermodynamics with generalized uncertainty principle Li Xiang, X. Q. Wen (Submitted on 6 Jan 2009 (v1), last revised 18 Mar 2009 (this version, v2))**

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| (189) In the standard viewpoint, the temperature of a stationary black hole is proportional to (=) its surface gravity, $T_H=\hbar\kappa/2\pi$. |
| (190)This is a semiclassical result and the quantum gravity effects are not (e) taken into consideration. |
| (191)This Letter explores a unified expression for (e) the black hole temperature in the sense of (e) a generalized uncertainty principle (GUP). |
| (192) Discussion involves a heuristic analysis of (e) a particle which is absorbed by (e) the black hole. |
| (193)Besides a class of static and spherically symmetric black holes, an axially symmetric Kerr-Newman black hole is considered. Different from the existing literature, authors suggest (eb) that the black hole's irreducible mass represent (eb) the characteristic size in (eb) the absorption process. |

(194)The information capacity of a remnant is also discussed by (e) Bousso's D-bound in de Sitter spacetime. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP 0910:046,2009 DOI: 10.1088/1126-6708/2009/10/046 Cite as: arXiv:0901.0603 [gr-qc] (or arXiv:0901.0603v2 [gr-qc] for this version)

IF we define

$$[a_7]_7 = [f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}]\,P_7^*$$

$$[a_7]_7 = \left[f_{7*}^{7}\right] P_7^* - \frac{h_*^7}{h_*} P^*(0)$$

$$e = [g_*^7]_7\, Q_7^*$$

$$\eta_7(t) = \left[h_*\right]_7 \beta_7\left(\theta_7(t)\right) - \mu_7(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\left[a_7\right]M_7 - [a_7]_7 W_7 - [n(t)]M_7$$

$$w_7(t) = eM_7$$

Write

$$A = \begin{pmatrix} -[a_7]_7 - [a_7]_7 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)= \begin{pmatrix} [\eta(t)]_7 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_7 \\ W_7 \end{pmatrix}$$

The solution of (32)is

$$X_7(t) = e^{A,t}[X]_7(0) + \int_0 e^{At}$$

## PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION

## ATTRITION MODELS

## MODULE EIGHT

Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal.

## 7. MODEL FORMULATION

**Model is applied to systems given in the following:**

**Researchers blur the line between classical and quantum physics by connecting chaos and entanglement July 12, 2016 by Sonia Fernandez**

(515)Experimental link between quantum entanglement (left) and (e&eb) classical chaos (right) found using (e) a small quantum computer. Credit: University of California - Santa Barbara Using a small quantum system

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| consisting of three superconducting qubits, researchers at UC Santa Barbara and Google have uncovered a link between aspects of classical and (e&eb) quantum physics thought to be unrelated: classical chaos and (e&eb) quantum entanglement. |
| (516)Their findings suggest (eb) that it would be possible to use (e) controllable quantum systems to investigate certain fundamental aspects of nature. |
| (517) "It's kind of surprising because (e) chaos is this totally classical concept—there's (=) no idea of chaos in a quantum system," Charles Neill, a researcher in the UCSB Department of Physics and lead author of a paper that appears in Nature Physics. |
| (518)"Similarly, there's no concept of entanglement within (eb) classical systems. And yet it turns out that chaos and (e&eb) entanglement are really very strongly and clearly related." |
| (519)Initiated in the 15th century, classical physics generally examines and describes (eb) systems larger than atoms and molecules. |
| (520) It consists of hundreds of years' worth of study including Newton's laws of motion, electrodynamics, relativity, thermodynamics as well as chaos theory—the field that studies the behavior of (e) highly sensitive and unpredictable systems. |
| (521) One classic example of chaos theory is the weather, in which a relatively small change in one part of (e) the system is enough to foil predictions—and vacation plans—anywhere on the globe. |
| (522)At smaller size and length scales in nature, however, such as those involving (e&eb) atoms and photons and their behaviors, classical physics falls short. |
| (523) In the early 20th century quantum physics emerged, with its seemingly counterintuitive and sometimes controversial science, including (e) the notions of superposition (the theory that a particle can be located in (eb) several places at once) and entanglement (particles that are deeply linked behave as (=) such despite physical distance from one another). |
| (524)And so began the continuing search for connections between the two fields. All systems are fundamentally quantum systems, according Neill, but the means of describing in a quantum sense the chaotic behavior of, (e) say, air molecules in an evacuated room, remains limited. |
| (525)Imagine taking a balloon full of air molecules, somehow tagging them so you could see them and then releasing them (eb) into a room with no air molecules, noted co-author and UCSB/Google researcher Pedram Roushan. |
| (526)One possible outcome is that the air molecules remain clumped together (e&eb) in a little cloud following the same trajectory around the room |
| (527) And yet, he continued, as we can probably intuit, the molecules will more likely take off in (eb) a variety of velocities and directions, bouncing off walls and interacting with (e&eb) each other, resting after the room is (=) sufficiently saturated with them. |
| (528)"The underlying physics is (=) chaos, essentially," he said. |
| (529)The molecules coming to rest—at least on (eb) the macroscopic level—is the result of (e) thermalization, or of reaching equilibrium after they have achieved (eb) uniform saturation within the system. |
| (530)But in the infinitesimal world of quantum physics, there is (=) still little to describe that behavior. |
| (531) The mathematics of quantum mechanics, Roushan said, do not allow for (e) the chaos described by Newtonian laws of motion. |
| (532)To investigate, the researchers devised an experiment using (e) three quantum bits, the basic computational units of (e) the quantum computer. |
| (533) Unlike classical computer bits, which utilize (e) a binary system of two possible states (e.g., zero/one), a qubit can also use (e) a superposition of both states (zero and one) as a single state. |

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| (534) Additionally, multiple qubits can entangle, or link so closely that their measurements will (=) automatically correlate. |
| (535)By manipulating these qubits with electronic pulses, Neill caused them to interact, rotate and evolve (e&eb, e&eb, eb) in the quantum analog of (e) a highly sensitive classical system. |
| (536)The result is a map of entanglement entropy of (e) a qubit that, over time, comes to strongly resemble that of (e) classical dynamics—the regions of entanglement in (eb) the quantum map resemble the regions of (e) chaos on the classical map. |
| (537)The islands of low entanglement in (eb) the quantum map are located in the places of (e) low chaos on the classical map. |
| (538)"There's a very clear connection between entanglement and (e&eb) chaos in these two pictures," said Neill. |
| (539) "And, it turns out that thermalization is (=) the thing that connects chaos and (e&eb) entanglement. |
| (540)It turns out that they are actually the driving forces behind (eb) thermalization. |
| (541)"What we realize is that in almost any quantum system, including on quantum computers, if you just let it evolve and you start to study what happens as a function of time, (eb) it's going to thermalize," added Neill, referring to the quantum-level equilibration. "And this really ties together the intuition between classical thermalization and (e&eb) chaos and how it occurs (eb) in quantum systems that entangle." |
| (542)The study's findings have fundamental implications for (e) quantum computing. |
| (543) At the level of three qubits, the computation is relatively simple, said Roushan, but as researchers push to build increasingly sophisticated and powerful quantum computers that incorporate (e) more qubits to study highly complex problems that are beyond the ability of classical computing—such as those in the realms of machine learning, artificial intelligence, fluid dynamics or chemistry—a quantum processor optimized for such calculations will be a very powerful tool. |
| (544)"It means we can study things that are completely impossible to study right now, once we get to bigger systems," said Neill.More information: Nature Physics, DOI: 10.1038/nphys3830 , http://www.nature.com/nphys/journal/vaop/ncurrent/full/nphys3830.html Journal reference: Nature Physics search and more info website Provided by: University of California - Santa Barbara |

**Title: coherent control of macroscopic quantum states in a single-Cooper-pair box Authors: Nakamura, Y.; Pashkin, Yu. A.; Tsai, J. S. Publication: Nature, Volume 398, Issue 6730, pp. 786-788 (1999). (Nature Homepage) Publication Date: 04/1999 Origin: NATURE DOI: 10.1038/19718 Bibliographic Code: 1999Natur.398...786N**

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| (545)A nanometre-scale superconducting electrode connected to (e&eb) a reservoir via (e&eb) a Josephson junction constitutes (e) an artificial two-level electronic system: a single-Cooper-pair box. |
| (546) The two levels consist of (e) charge states (differing by 2e, where (e) e is the electronic charge) that are coupled by (e) tunnelling of Cooper pairs through (e&eb) the junction. |
| (547)Although the two-level system is macroscopic, containing (e) a large number of electrons, the two charge states can be (=) coherently superposed. |
| (548)The Cooper-pair box has therefore been suggested as (=) a candidate for a quantum bit or `qubit'-the basic component of a quantum computer. |
| (549)Here authors report the observation of quantum oscillations in (eb) a single-Cooper-pair box. |

(550)By applying a short voltage pulse via (e&eb) a gate electrode, we can control (e&eb) the coherent quantum state evolution: the pulse modifies (e&eb) the energies of (e) the two charge states non-adiabatically, bringing them into (e&eb) resonance.

(551)The resulting state-a superposition of (e) the two charge states-is detected by (e) a tunnelling current through (e&eb) a probe junction.

(552)Results demonstrate electrical coherent control of (e) a qubit in a solid-state electronic device.

**Superconducting quantum computing From Wikipedia, the free encyclopedia**

(553)Superconducting quantum computing is a promising implementation of (e) quantum information technology that involves (e&eb) nanofabricated superconducting electrodes coupled through (e&eb) Josephson junctions.

(554)As in a superconducting electrode, the phase and (e&eb) the charge are conjugate variables.

(555) There exist three families of superconducting qubits, depending on (=) whether the charge, (e&eb) the phase, or neither of (e) the two is good quantum number.

(556) These are respectively termed charge qubits (e&eb) flux qubits, and (e&eb) hybrid qubits.

**Theory**

(557)Unlike many other physical implementations of a qubit which involve (e&eb) exclusively two-level systems (such as nuclear spin and (e&eb) photon polarization), the integrated quantum circuit involved in (e&eb) a superconducting qubit is (=) a multi-level system, of which only the first two levels are used as (=) the computational basis.

(558)A basic requirement for such an implementation is that the energy levels are (=) not uniformly spaced, so that (e) photons of a particular frequency which cause (eb) transition between the 0 and 1 levels do not(e) cause transitions from the first level to (e&eb) the higher levels as well.

(559)For electronic signals to be carried from one part of the circuit to (e&eb) another without energy loss and (e&eb) hence decoherence, the system also needs to be (=) non-dissipative, i.e., the metallic parts involved should have (e) zero resistance.

(560) These circuits currently need to be (=) operated at very low temperatures so that, 1) superconductivity is (=) realized, and 2) thermal fluctuations do not (e) cause transitions between (e&eb) energy levels.

**Quantum LC Circuit**

(561)The simplest non-dissipative quantum circuit consists simply of (e) an inductor L {\displaystyle L} L and (e&eb) capacitor C {\displaystyle C} C, with (e&eb) the metallic wires connecting them being superconducting.

(562)Here the flux Φ {\displaystyle \Phi } \Phi through (e&eb) the inductor and (e&eb) the charge Q {\displaystyle Q} Q in the capacitor are (=) canonically conjugate variables which obey (eb) the commutation relation [ Φ , Q ] = i ℏ {\displaystyle [\Phi ,Q]=i\hbar } [\Phi ,Q]=i\hbar

**The Hamiltonian associated with this system is**

(563)H = Φ 2 2 L + Q 2 2 C {\displaystyle H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}} H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}

(564)This gives rise to the energy levels E = ħ ω 0 ( n + 1 2 ) {\displaystyle E=\hbar \omega _{0}(n+{\frac {1}{2}})} E=\hbar \omega _{0}(n+{\frac {1}{2}}) Here ω 0 = 1 L C {\displaystyle \omega _{0}={\frac {1}{\sqrt {LC}}}} \omega _{0}={\frac {1}{{\sqrt {LC}}}}

**Quantum gravity Carlo Rovelli (2008), Scholarpedia, 3(5):7117. doi:10.4249/scholarpedia.7117 revision #135539 [link to/cite this article] Post-publication activity Curator: Carlo Rovelli Dr. Carlo Rovelli, Center of Theoretical Physics, and Marseille, France**

(565)Quantum Gravity is the name given to any theory that describes gravity in (eb) the regimes where quantum effects cannot be disregarded.

(566)At present, there is no such a theory which is universally accepted and confirmed by experience. Therefore the term "Quantum Gravity" indicates (eb) more an open problem than a specific theory.

(567)Several research lines, at different levels of development, offer tentative solutions to the problem. These tentative quantum-gravity theories are variously viewed as competing research directions, or as contributions to the common goal of finding the physically correct theory. The quest for the good quantum theory of gravity bears on a number of fundamental issues and it is sometimes presented as the most important open problem in fundamental physics: the "Holy Grail" of contemporary theoretical physics.

**The problem of quantum gravity**

(568)The current theory of (e) gravity is general relativity, a theory developed by Albert Einstein at the beginning of the XX century, which has changed in depth our understanding of (e) what are space and time.

(569) Similarly, quantum mechanics has changed in depth our understanding of matter, energy and causality. These two theories, today very well empirically supported, have opened (eb) a major conceptual revolution in physics.

(570)However, they appear to be incompatible, at least at a first reading, because (e) each of the two is formulated on the basis of (e) principles that are explicitly contradicted by (e) the other theory.

(571) Therefore in spite of the immense development of scientific knowledge it has delivered, the physics of the XX century has left us with (e&eb) a major confusion on what are the basic conceptual ingredients for (e) understanding the physical world.

(572)In other words, the scientific revolution opened by general relativity and quantum mechanics at the beginning of (e) the XX century is not concluded yet, and a new synthesis is required. Quantum gravity, merging general relativity and quantum mechanics should be this new synthesis.

**Quantum space and quantum time**

(573)Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.

(574) In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.

(575) They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.

(576)When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.

(577) Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.

**The Planck scale**

(578)Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".

(579)Here $\hbar$ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.

(580)The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.

(581)Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)

(582)Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.

**Physical applications**

Some phenomena where quantum gravity plays a major role are the following.

**The microscopic structure of spacetime**

(583)If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.

(584) Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length

(585) Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.

(586) A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.

**Early cosmology**

(587)According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.

(588) However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.

(589) A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.

(590) Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".

(591)One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.

**Black holes**

(592)Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.

(593) First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.

(594) Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.

(595)What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.

(596) Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.

**Astrophysical effects**

(597)Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.

(598) An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.

(599)The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.

**Main approaches**

(600)The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct

bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.

(601) Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.

(602)The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.

**String theory**

(603)String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.

(604) Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.

(605)String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.

(606)Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.

(607)However, when summed the entire series appears to be (=) divergent.

(608)In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.

(609)Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.

(610)Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.

(611) In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.

**Loop quantum gravity and spin foams**

(612)Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.

(613)The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity.

(614)Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime.

(615)The granular properties of space can be (=) explicitly computed.

(616) In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory.

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| (617)The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks. |
| (618) Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure. |
| (619) The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit. |
| (620)Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism. |
| (621)This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism. |

**Noncommutative geometry**

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| (622)Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime. |
| (623) Quantum mechanics teaches us that dynamical quantities are (=) noncommutative --meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed. |
| (624)It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry. |
| (625) A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces. |
| (626)Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces. |
| (627)This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity. |

**Others**

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| (628) Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays |
| (629). Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection. |
| (630) Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible. |

(631) The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.

Let $p_8(t, a_8)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(g_8)Q_8$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_8(0) = 0, \quad g_8(P_8) > 0, \quad P_8 \text{ 0}, \qquad \lim_{P_8 \to \infty} g_8(P_8) = (K_8)_8 \text{ }\infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_8 + (\rho_a)_8 + D_8(a_8, P_8, Q_8)\rho_8 + 0, \qquad 0 < a_8 \text{ }\infty,$$

$$\rho_8(t,0) = \int_0^\infty n_8(a_8, P_8, Q_8)\rho_8(t_8, a_8)da_8$$

Where $D_8$ is the mortality function and $n_8$ is the fertility function.

Put $_8$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_8}{dt}\right]_8 = x_8\left[\frac{dQ_8}{dt}\right]_8 - x_8\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_8}{dt}\right]_8 = x_3\left[\frac{dQ_3}{dt}\right]_8 - x_4\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_8}{dt}\right]_8 - x_6\left[\frac{dQ_3}{dt}\right]_8$$

Where $x_8, x_8, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_8$ and $g_8(P_8)Q_8$

Now when you add the term $(b_8Q_8)$ with $g_8(P_8)(Q_8)$ we are infect see is that $(b_8Q_8)$ is credited incrementally part of $g_8(P_8)(Q_8)$ or instant nearly, when subtracted from $\frac{dQ_8}{dt}$ , it means that $g_8(P_8)(Q_8)$ is slowly eaten by $\frac{dQ_8}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_8 = \mu_8(a_8) + (f_8)_8(P_8)_8 + (f_8)_8(Q)_8 \quad (2)$$

$$n_8 = \beta_8(a_8)h_8(Q_8), \quad (3)$$

Where $\mu_8$the age- dependent debit rate whit is $\mu_8(a_8)$ 0 $\forall a > 0$; $(f_8)_8$ is the density- dependent death rate with $f_8(0) = 0$, $(f_8^8)_8(P_8)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_8)_8(P_8) = (k_8)_8$ ; $f_8$ is the functional response of creditors to prey with $[f_8]_8(0) = 0, (f_8^8)_8(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_8)_8(Q_8) = [k_8]$ ; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^8(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_8 + (\rho_a)_8 = -\mu_8(a_8) + (f_8)_8P_8) + \big((f_8)_8(Q_8)\big)(\rho_8)_8 \quad (4)$$

$$\rho_8(t_8, 0) = h_8(Q_8) \int_0^{\infty} \beta_8(a_8)\rho_8(t, a_8)da, \quad (5)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(P_8)Q_8 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_8$ and $Q_8$ of (4-6) satisfy

$$0 \leq \rho_8(t, a_8) \leq \rho_8^*(t, a_8),$$

$$0 \leq Q_8(t) \leq Q_8^*(t),$$

Where $\rho_8^*(t, a_8)$ and $Q_8^*(t)$, are solutions of

$$(\rho_t^*)_8 + (\rho_a^*)_8 = -(\mu_8(a_8) + (f_8)_8(P_8^*) + (f_8)_8 Q_8^*))\,\rho_8^*, \quad (7)$$

$$\rho_8^*(t, 0) = \int_0^{\infty} \beta_8(a_8)\,\rho_8^*(t, a_8)da_8 \quad (8)$$

$$\frac{dQ_8^*}{dt} = -b_8 Q_8^* + g_8(P_8^*)Q_8^*, \quad (9)$$

With $P_8^*(t) = \int_0^{\infty} \rho_8^*(t, a_8)da_8$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_8^*(t)}{dt} = P_8^*(t)\big(r_8 - (f_8)(P_8^*) - (f_8)_8(Q^*)\big), \quad (10)$$

$$\frac{dQ_8^*(t)}{dt} = Q_8^*(t)(-b_8 + g_8(P_8^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_8(a)e^{-\int_0^a \mu_8(\tau)dt}e^{-ra}\,da_8 = 1. \qquad (12)$$

If $r < 0$, the trivial solution $P_8^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(t)dt}\,da_8 > 1, \qquad (H1)$$

Which guarantees $r_8 > 0$.

*Remark.* If $\beta_8(a) = [\beta_0]_8 a_8 e^{-\alpha_8 a_8}, \mu_8(a_8) \equiv \mu_8$ then the solution of (12) is

$$r_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_8]_8(P_8) = 0, [f_8](Q_8) = \lambda_8 Q_8\ h_8(Q_8) = 1, \qquad and\ \ g_8(P_8) = c_8 P_8$$

$$[\rho_8]_8 + [\rho_a]_8 = -(\mu_8 + \lambda_8 Q_8)\rho_8$$

$$\rho_8(t,0) = \int_0^\infty [\beta_0]_8 a_8 e^{-\alpha_8 a_8}\rho_8(t,a_8)\,da_8$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + c_8 P_8 Q_8$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_8}{dt} = \omega_8 P_8 - \lambda_8 P_8 Q_8$$

$$\frac{dQ_8}{dt} = b_8 Q_8 - c_8 P_8 Q_8$$

Where $\omega_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$

## 8. THE EQUILIBRIUM $(\rho_0)_8(a_8), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_8(a_8), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_8 + (\rho_a)_8 = -\big(\mu_8(a_8) + (f_8)_8(P_8)\big)\rho,$$

$$\rho_8(t,0) = \int_0^\infty \beta_8(a_8)\rho_8(t,a_8)da_8 \qquad (13)$$

The solution of (13) has the form of

$$\rho_8(t,a_8) = Ce_8^{r(t-a_8)} - \int_0^a (\tau)d\tau - \int_0^8 (f_8)_8\big(P_8(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_8(t) = C_8 e^{r_8 t - \int_0^t (f_8)_8(P_8(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_8^{-8}(r)$ if $k_8 > r$. This is equivalent to

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{(k_8)_8 a_8} da_8 < 1.$$

In fact, if we define

$$R_8(s_8) = \int_0^\infty \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau)d\tau} e^{-s_8 a_8}\, da_8$$

$R(s)$ is a decreasing function. $R_8(r_8) = 1$ And $R_8(k_8)_8 < 1$ imply $r_8 < (k_8)_8$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_8(a_8), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{k_8 a_8} da_8$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_8 \equiv \int_0^\infty (\rho_0)_{88}(a_8)da_8$ *satisfies*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{-(f_8)_8 (P_0)a_8} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_8(a_8), 0)$

Let $u_8(t, a_8) = \rho(t, a_8) - \rho_0(a_8)$ Then

$$N_8(t) \equiv \int_0^\infty u_8\,(t, a_8)da_8 = \int_0^\infty \rho_8(t, a_8)da_8 - \int_0^\infty (\rho_0)_8(a_8)da = P_8(t) - P_0)_8$$

And

$$(u_8)_8 + (u_a)_8 = (\rho_t)_8 + (\rho_a)_8 - (\rho_0^8)_8(a_8)$$

$$= -\mu_8(a_8) + (f_8)_8(P_8) + (f_8)_8(Q_8))\,(u_8 + (\rho_0)_8) + \mu_8(a_8)(\rho_0)_8 + (f_8)_8(P_0)_8(\rho_0)_8$$

$$= -\mu_8(a_8)u_8 - (f_8)_8(P_8) + (f_8)_8(Q_8))\,u_8 - (\rho_0)_8\,(f_8)_8(P_8) - (f_8)_8(P_0)_8$$

$$+ (f_8)_8(Q_8)), \qquad (15)$$

Where $(\rho_0^8)_8(a_8)$ denotes $d(\rho_0)_8(a_8)/da_8$

Expansion $(f_8)_8(P_8)$ and $(f_8)_8(Q_8)$ at $(P_0)_8$ and 0 yield respectively

$$(f_8)_8(P_8) = (f_8)_8((P_0)_8 + N_8)$$

$$= (f_8)_8(P_0)_8 + (f_8^8)_8(P_0)_8N_8 + o_8(N_8), \qquad (16)$$

$$(f_8)_8(Q_8)$$

$$= (f_8^8)_8(0)Q + o_8(Q_8) \qquad (17)$$

Write $f_8(P_0) = f_{80}, f_8'(P_0) = f_{80}'$, and $f_8'(0) = f_{80}'$ and substitute (16) and (17) into (15) to obtain

$$(u_8)_8 + (u_a)_8 \approx -(\mu_8(a_8) + (f_{80})_8)u - (f_{80}^8)_8(\rho_0)_8(N_8) - (f_{80}^8)_8(\rho_0)_8(Q_8)$$

$$u_8(t,0) = \rho_8(t,0)(\rho_0)_8(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a) + (\rho_0)_8(a_8)da_8 - (\rho_0)_8\,(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a_8)da_8 + (h_8(Q_8) - 1)(\rho_0)_8\,(0) \qquad (18)$$

Expanding

$$h_8(Q_8) = 1 + (h^8)_8(0)Q_8 + 0(Q_8) \equiv 1 + h_0'Q_8 + 0Q_8$$

And substituting it into (18), we arrive at

$$u_8(t,0) \approx \int_0^\infty \beta_8\,(a_8)u_8(t,a_8)da_8 + (h_{a_8}')(\rho_0)_8Q_8$$

The linearization is the following :

$$[u_8] + [u_8] = -(\mu_8(a_8) + [t_{80}]_8 u_8 - [\alpha_8]_8(a_8)N_8 - [\alpha_8]_8(a_8)Q_8$$

$$u_8(t,0) = \int_0^\infty \beta_8(a_8)u_8(t,a_8)da_8 + [\xi_8]Q_8 \qquad (19)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8$$

where

$$[\alpha_8]_8(a_8) = [t_8^8 a]_8[p,o]_8(a_8)d2(a_8) = [f_{80}^8]_8(\rho\, o\,)_8(a_8)\xi = [\mathrm{h\ o}^8]_8[\rho \mathrm{o}]_8(\mathrm{o}) \text{ and } \mathrm{b}_8 = \mathrm{b}_8 - g_8(P_8)$$

If $g_8(P_0) > b$, clearly, $(\rho_0)_8(a_8), 0)$ is unstable and the predator population is persistent. Assume $g_8(P_0)_8 < b$.

Then, $Q_8(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_8 + u_a = -(\mu_8(a_8) + (f_{80})_8)u_8 - (\alpha_8)_8(a_8)N_8$$

$$= -(\mu_8(a_8) + f_{80})u_8 - \alpha_8(a_8)\int_0^\infty u_8(t,a_8)da_8$$

$$u_8(t,0) = \int_0^\infty \beta_8\,(a_8)u_8(t,a_8)da_8$$

Using the method developed by Cushing [2, 3], we suppose

$$u_8(t,a_8) = \tilde{\mathrm{u}}_8(a_8)^{e^{\gamma(t-a_8)}}$$

And define $w = \int_0^\infty \tilde{u}(a) e^{-\gamma a_8}\, da_8$. Then,

A solution of the first differential equation for $\tilde{u}(a_8)$ gives,

$$\tilde{u}_8(a_8) = e^{-\int_0^a \mu_8(\tau_8)d\tau - (f_{80})_8 a}\left(\tilde{u}_8(0) - w\int_0^a e^{\int_0^\tau \mu_8(s_8)ds + (f_{80})}\tau_8\,(\alpha_8)_8(\tau_8)e^{\gamma_8\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\right)\tilde{u}_8(0)$$

$$+w\int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\int_0^a (\alpha_8)_8(\tau_8)e^{-\gamma_8\tau_8}e^{(f_{80})_8\tau_8 + \int_0^\tau \mu_8(s_8)ds}d\tau_8\, da_8$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}da\right)(\tilde{u}_8(0))_8$$

$$+\left(-1 - \int_0^\infty e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\int_0^a e^{(f_{80})_8\tau_8 + \int_0^\tau \mu_8(s_8)ds + y\tau}\alpha_8(\tau)d\tau\, da\right)w = 0.$$

(23)

Define $Y(\gamma, a) = e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(t_8)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_8\,(a_8)Y_8(\gamma_8,a_8)da_8\right)u_8(0)$$

$$+w\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)\int_0^{a}(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8=0,$$

(24)

$$\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\tilde{u}_8(0)$$

$$+\left(-1-\int_0^{\infty}Y_8(\gamma_8,a_8)\int_0^{a}(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8\right)w=0.$$

Equation (20) has a solution of the form of $u_8(t,a_8)=u_8(a_8)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_8(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_8(\gamma_8,a_8)=\int_0^{a}(\rho_0)_8(\tau_8)Y^{-8}\,(\gamma_8,a_8)d\tau_8$. We obtain the characteristic equation

$$f_{80_8}^{l}\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\left(\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$+\left(1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)da\right)\left(1+(f_{80}^{l})_8\int_0^{\infty}Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_8(a_8)=(\rho_0)_8(0)e^{-\int_0^{a}\mu_8(\tau_8)d\tau_8-(f_{80})_8a_8}$$

$$I_8(\gamma_8, a_8) = \int_0^a (\rho_0)_8(\tau_8) e^{\gamma_8\tau_8 + (f_{80})_8\tau + \int_0^\tau \mu_8(s_8)ds_8} d\tau$$

$$= \frac{(\rho_0)_8(0)}{\gamma_8}(e^{\gamma_8 a_8} - 1)$$

Let

$$E_8(a_8) = e^{-(f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau_8}$$

Then $[Y_8]_8[\gamma_8]_8, a_8) = E_8(a_8)e^{-\gamma_8 a_8}$. by a simple calculation,

$$\int_0^\infty \beta_8(a_8)Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma}\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}\ da\right);$$

$$\int_0^\infty Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma_8}\left(\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}\ da_8\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_8(0)}{\gamma_8}[f_{80}^8]\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}da_8\ (\int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8)$$

$$+\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right).$$

$$\left(1 + \frac{(\rho_0)_8(0)}{\gamma_8}(f_{80}^{'})_8\left(\int_0^\infty E_8(\,a_8)da_8 - \int_0^\infty E_8(\,a_8)e^{-\gamma_8 a_8}da_8\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right)(1+\frac{(f_{80}^{'})_8}{\gamma_8}(\rho_0)_8(0)\int_0^{\infty}E_8(a_8)da_8)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)da_8=1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\neq 0,$$

Since $\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)\int_0^{a}\alpha_8(\tau_8)Y^{-8}(\gamma_8,\tau_8)d\tau_8 da_8>0$. Therefore, the characteristic solution is

$$\gamma_8=-(f_{80}^{'})_8(\rho_0)_8\int_0^{\infty}E_8(a_8)da_8.$$

Since $f_8'(P_0)_8>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_8),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_8),0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a),Q^*)$. It must satisfy the equations:

$$[p_a^*]_8 + (\mu_8(a_8) + [f_8]_8(P_8^*) + [f_8]_8(Q_8^*))P_8^* = 0, \quad (26)$$

$$p_8^*(0) = h_8(Q_8^*)\int_0^\infty \beta_8(a_8)p_8^*(a_J)p_8^*(a_8)da_8 \quad (27)$$

$$g_8(P_8^*) - b_8 = 0, \quad (28)$$

Where $P_8^*=\int_0^\infty p_8^*(a_8)\,da_8$

Equation (28) yields $P^* = g^{-8}(b)$,

From (26),

$$p_8^*(a) = p_8^*(0)e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}$$

Substituting into (27), we need to solve

$$\left(1 - h_8(Q_8^*)\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}da_8\right)p_8^*(0) = 0$$

For $Q_8^*$

Fix $p_8^*(0)$ and $P_8^*$ Let

$$H_8(z_8) = 1 - h_8(z_8)\int_0^\infty \beta_8(a_8)\,e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}da_{,8}$$

And assume that

$$[k_3]_8\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(g_8^{-8}(b_8))a_8 - [k_8]_8a_8}\,da_8 < 1. \quad (H3)$$

Then

$$\lim_{z\to\infty}[H_8(z_8)] > 0.$$

Clearly,

$$\frac{dH_8}{dz_8} = h_8(z)([f_8^{'}]_8(z_8))\int_0^{\infty} a_8\beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - f_8(P_8^*)a_8 - f_8(z_8)a_8}\, da_8$$

$-h^{'}(z_8)\int_0^{\infty}\beta_8(a_8)\, e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}\quad da_8$

By the assumptions on $\left[f_8^8\right]_8 and\ [h^8]^8 dH_8/dz > 0$. Hence, if

$$H_8(0) = 1 - \int_0^{\infty}\beta_8(a_8)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8}da_8 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_8(0) < 0$ follows if $g_8^{-}(b_8) < [f_8^{-8}]_8(r_8)$, where $r_8$ is the solution of (12).

Indeed, since

$$R_8(s_8) = \int_0^{\infty}\beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 e^{-y_8^a}}da_8$$

Is a decreasing function of $s_8$ when $P_8^* = [g^{-8}]_8(b_8) < [f_8^{-8}]_8(r_8)$, that is $[f_8]_8(P_8^*) < r$, $R\left(f_8(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_8(g^{-8}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8\left(g_8^{-8}(b)\right)a_8} > 1. \quad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_8^*(a_8), Q_8^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_8^*(a_8), Q_8^*)$ is more complicated. We again consider perturbations $m_8(t_8 a_8) = p_8(t_8, a_8) - p_8^*(a_8)$ and $W_8(t) = Q_8(t) - Q^*$.Then

$$M_8(t) \equiv \int_0^\infty m_8(t, a_8) da_8 = P_8(t) - P_8^*,$$

$$[m_8]_8 + [m_a]_8 = [p_8]_8 + [p_a]_8 - \left[p^*\right]_8(a_8) \approx -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8$$

$$-[f_{8*}^8]_8 p_8^*(a_8)M - [f_{8*}^8]_8 p_8^*(\mathrm{a})W_8,$$

$$m_8(t,0)p_8(t,0) - p_8^*(0) \approx h_* \int_0^\infty \beta_8(a_8) m_8(t, a_8)\, da + \frac{[h_*^8]_8}{[h_*]_8}\, p_8^*(0) W_8$$

$$\frac{dW_8}{dt} = \frac{dQ_8}{dt} \approx [g_*^8]_8 Q_8^* M_{,8}$$

Where $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $\left[f_{8*}^{'}\right]_8$ denoting $\left[f^{'}{}_8(P^*)\right]_8 \left[f_{8*}^{'}\right]_8$ denoting $[f_8]_8$ $(Q_8^*,$ $h_{*}^{'}$ denoting $h_{*}^{'}$ denoting $h_8^{'}$ $(Q_8^*)$, and $\left[g_*^{'}\right]_8$

Denoting $\left[g^{'}\right]_8(P_8^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_8]_8 + [m_a]_8 = -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8 - [f'_{8*}]_8[p_8^*](a_8)M_8 - [f_{8*}^{'}]_8 p_8^*(a_8)W_8$$

$$m_8(t,0) = [h_*]_8 \int_0^\infty \beta_8(a_8)m_8(t,a_8)da_8 + \frac{[h_*^8]_8}{[h^*]_8}[p^*]_8(0)\, W_{,8} \qquad (30)$$

$$\frac{dW_8}{dt} = [g_*^8]_8[Q_8^*]M_8$$

Integrating $[m_8]_8(t,a_8)$ with respect to a from 0 to ∞, we have

$$\frac{dM_8(t)}{dt} = \int_0^\infty [mt]_8(t,a_8)da_8$$

$$= -\int_0^\infty [m_a]_8\,(t,a_8)da_8 - \int_0^\infty \mu(a_8)m(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + [f'_{8*}]_8 \int_0^\infty p_8^*(a_8)da_8\right)M_8 - [f_{8*}^8]_8 \int_0^\infty p_8^*(a_8)da_8 W_8$$

$$= m_8(t,0) - \int_0^\infty \mu_8(a_8)\, m_8(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + [f_{8*}^{'}]_8 P_8^*\right)M_8 - [f_{8*}^{'}]_8 P_8^* W_{,8}$$

From (30), we arrive at the equations

$$\frac{dW_8(t)}{dt} = \int_0^\infty ([h_*]_8\beta_8(a_8) - \mu_8(a_8))m_8(t,a_8)da_8$$

$$-\left([f_{8*}] + [f_{8*}]_8 + [f_{8*}^{'}]_8 P_8^*\right)M_8 + \left(\frac{[h_*^{'}]_8}{[h_*]^8}\, P_8^*(0) - f_{8*}^{'} P_8^*\right)W_8, \qquad (31)$$

$$\frac{dW_8(t)}{dt}\left[g^{'}{}_{*}\right]_8 Q_8^* M_{.8}$$

By the mean value theorem,

$$\int_0^{\infty}([h_{*}]_8\beta_8(a_8)-\mu_8(a_8))m_8(t,a_8)da_8=([h_{*}]_8\beta_8(\theta_8)-\mu_8(\theta_8))M_{,8}$$

Where $0<\ \theta_8=\theta_8(t)<\infty$.

And if $a_8^8[a_8^8]2<4\ [a_{\ 8}]2<4\ a_{\ 8}e$,

$\mathrm{Re}\,\lambda_{\ 8}\left[\lambda_{\ 8}\right]_{\ 8}=\mathrm{Re}\lambda_{\ 8}\left[\lambda_{\ 8}\right]_{\ 8}=-\frac{8}{8}a_8[a2\ ]_{\ 8}<0.$

Let $\Delta=\max_{8<u<\infty}\left|h_{\ *}\,[h_{\ *}]_{\ 8}\beta_{\ 8}(a_8)-\mu_8(a_8)\right|$. Then$|n_8(t)|$ 2∆∀t 0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^{*}(a),\boldsymbol{Q}^{*})$ is asymptotically stable if either $a_8^8<4\ a_{\ 8}\ e$ and $[a_{\ 8}]>1\Delta$,or $[a_8^8]$ 24$[a_{\ 8}]e$ and $(a_{\ 8}-\sqrt{(a_8^8)_{\ 8}-4(a_{\ 8})_{\ 8}\,e})>1\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a)=\mu_0+\left(\frac{a}{a+8}\right)$ and $\beta_{\ 8}(a_8)=ae^{-ma}$ Then $\Delta=[\mu_0]+2$ Simple calculation leads to

$$\int_0^{\infty}\beta_{\ 8}(a_8)e^{-\int_0^a\mu_8(\tau_{\ 8})+d\tau_8 da}=\frac{1}{(m+\mu_0+1)^{\ 8}}+\frac{1}{(m+\mu_0+1)^{\ 3}}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_8 + [\mu_0]$ 0.5. Assume$[f_8](x_8)= K_8x_8$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_8(x_8)= [0x_8]/(2 + x_8)$; the functional responses $[[f_8](x) = [0]\, x_8/(1 + x_8)$and $h_8(x_8)$=(1+$\omega_8x_8$)$/(1 + x_8)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_8 + [\mu_0] + [\mathrm{K}][P_0])}{[m_8]_{\,8} + (\mu_0)_{\,8} + [\mathrm{K}][P_0])\,{}_8^3\ [\mathrm{K}]_{\,8}[P_0]_8^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_8$i e $[P_0] < b/(\delta_8 - b_8)$then the equilibrium$(P_0)$ $[P_0]_{\,8}(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_8$

$[f_{8*}] = K_8\ \mathrm{P}^*\ [f_{8*}] = \frac{\sigma \mathrm{Q}_8^*}{8+\mathrm{Q}_8^*}$ $[h_*] = \frac{8+\omega_8 \mathrm{Q}_8^*}{8+\mathrm{Q}_8^*}$

$[f_{8*}] = K_8\ \left[\mathrm{f}_{8*}^{\,8}\right] = \frac{\sigma_8}{(8+\mathrm{Q}_8^*)^8}[\ \mathrm{g}_*^{\,8}] = \frac{\sigma_8}{(8+\mathrm{P}_8^*)^8}$ $\left[\mathrm{h}_*^{\,8}\right] = \frac{\omega_8 - 8}{(8+\mathrm{Q}_8^*)^8}$

In order to have (H3) and (H4) hold,$m, \mu_0,$ K, b, δ, and $\omega$ need to satisfy

$$\frac{\omega_8}{(m + [\mu_0]_{\,8} + 1 + \frac{K_8 b_8}{\delta_8 + b_8} + \sigma)^{\,8}} + \frac{1\omega_8}{([m_8]_{\,8} + [\mu_0]_{\,8} + 1 + \frac{K_8 b_8}{\delta_8 + b_8} + \sigma_8)^{\,3}} < 1$$

And

$$\frac{1}{(m_8 + [\mu_0]_{\,8} + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^{\,8}} + \frac{1\omega_8}{(m_8 + [\mu_0]_{\,8} + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^{\,3}} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8} + (m_8 + [\mu_0]_8 + \frac{K_8 b_8}{\delta_8 - b_8} + \sigma_8)\, Q_8^*)^3$$

$$= m_8 + [\mu_0]_8 + 3 + \frac{K_8 b_8}{\delta_8 - b_8}(1 + Q^*)(1 + \omega_8 Q_8^*$$

Now

$$\left[a_1^1\right] - 4\left[a_8\right]e = \left(\frac{1K_8 b_8}{\delta_8 - b_8}\right)^8 + \frac{4K_8 b_8 \sigma_8 Q_8^*}{(\delta_8 - b_8)(8 + Q_8^*)} + \frac{\sigma_8^8 Q_8^{*8}}{(1 + Q_8^*)^8}$$

$$\frac{4\, b_8 \sigma_8 (\delta_8 - b)Q^*}{\delta_8 (1 + Q_8^*)^8} - \frac{4\,(\delta_8 - b_8)^8 (1 - \omega_8) Q_8^*}{\delta_8 (1 + \omega_8 Q_8^*)(1 + Q_8^*)} P * (0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_8\, ae_8^{mu}\ \ [\mu_0]_8 - \frac{a_8}{a_8 + 8}$$

Hence if $[a_8^8 > 4(a_8)_8 e$ and

$$[a_8] = \frac{1K_8 b_8}{\delta_8 - b_8} + \frac{\sigma_8 Q_8^*}{1 + Q_8^*} > 1\Delta$$

**The nontrivial equilibrium is stable.**

IF we define

$$[a_8]_8 = [f_{8*}]_8 + [f_{8*}]_8 + [f_{8*}]\,P^*_8$$

$$[a_8]_8 = \left[f^{8}_{8*}\right]P^*_8 - \frac{h^8_*}{h_*}P^*(0)$$

$$e=[g^8_*]_8\,Q^*_8$$

$$\eta_8(t) = \left[h_*\right]_8\,\beta_8\left(\theta_8(t)\right) - \mu_8\left(\theta(t)\right),$$

The system(32)can be written as the following system:

$$[M(t)]\left[a_8\right]M_8 - [a_8]_8 W_8 - [n(t)]M_8$$

$$w_8(t) = eM_8$$

Write

$$A=\begin{pmatrix} -[a_8]_8 - [a_8]_8 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_8 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_8 \\ W_8 \end{pmatrix}$$

The solution of (32)is

$$X_8(t) = e^{A,t}\,[X]_8(0) + \int_0 e^{At}$$

This equation notwithstanding its importance in all the models stand out for its importance for the following models

to which it is applied:

**Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016**

(111)These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field.

(112)The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets.

(113) Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state.

(114) Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands.

(115)The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI

(116) The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9].

(117)This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive.

(118)More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

(119)Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.

(120)Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.

(121) The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.

(122)This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.

(123)The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.

(124)This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.

(125) For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.

(126)The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(127)Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.

(128)We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.

(129)When this is true—as is often the case—the system serves as (=) a “bath” for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called **thermalization.**

(130)However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.

(131)Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1).

(132)Bringing them to thermal equilibrium requires (e) an external bath.

| |
|---|
| (133) But so far, there has been no rigorous proof of (e) the existence of **such many-body localization**. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption. |
| (134)The concept of **localization** was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations). |
| (135)The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century. |
| (136)But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information. |
| (137)Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment. |
| (138)For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood. |
| (139) It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7]. |
| (140) Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature. |
| (141) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators. |
| (142)However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8]. |
| (143)Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems. |
| (144) Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL. |
| (145)As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked. |
| (146)His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice. |
| (147) An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site. |
| (148)This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself). |

(149)This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition.

(150)However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime:

(151) These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects.

(152)It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies).

(153) This is a realistic scenario for (e) all known systems.

(154)On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects.

(155)On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other.

(156)No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true.

(157) His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found.

(158)Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization?

(159)Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects?

(160)Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions.

(161)So in those systems, the quantum phase transition does indeed exist.

(162)But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension.

(163)Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters.

**PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION ATTRITION**

**MODELS**

**MODULE NINE**

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Model is applied to following formulations Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. 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**(e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP This paper is divided into four parts. In part A, some general considerations about (e&eb) gravitational radiation are followed by (e) a treatment of the scalar wave equation in the manner later to be applied to (e&eb) Einstein's field equations. In part B, a co-ordinate system is specified which is suitable for (e) investigation of outgoing gravitational waves from (e) an isolated axi-symmetric reflexion-symmetric system. The metric is expanded in (eb) negative powers of (e) a suitably defined radial co-ordinate r, and the vacuum field equations are investigated in detail. It is shown that the flow of information to infinity is controlled by (e&eb) a single function of (e) two variables called the news function. Together with initial conditions specified on a light cone, this function fully defines (eb) the behaviour of the system. No constraints of (e) any kind are encountered. In part C, the transformations leaving (eb) the metric in the chosen form are determined. An investigation of the corresponding transformations in (eb) Minkowski space suggests (eb) that no generality is lost by (e) assuming that the transformations, like the metric, may be expanded in (eb) negative powers of r. Gravitational Waves in General Relativity VII Waves from Axi-Symmetric Isolated Systems Royal Society Proceedings H. Bondi, M. G. J. van der Burg, A. W. K. Metzner Published 21 August 1962.DOI: 10.1098/rspa.1962.0161 STOP There is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant $\Lambda$, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive $\Lambda$ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the $\Lambda>0$ generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible. Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of (e) GW signals from BBHs with component spins $\chi_{1,2}\in[-0.99,0.99]$ aligned (e&eb) with the orbital angular momentum, component masses $m_{1,2}\in[2,48]M_\odot$, and (e&eb) total mass $M_{total}\leq 50\ M_\odot$. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins $>+0.6$ and component masses $\gtrsim 20\,M_\odot$. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201**

## 9. MODEL FORMULATION

Let $p_9(t, a_9)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(g_9)Q_9$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_9(0) = 0, \quad g_9(P_9) > 0, \; P_9 \; 0, \qquad \lim_{P_9 \to \infty} g_9(P_9) = (K_9)_9 \; \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_9 + (\rho_a)_9 + D_9(a_9, P_9, Q_9)\rho_9 + 0, \qquad 0 < a_9 \; \infty,$$

$$\rho_9(t,0) = \int_0^\infty n_9(a_9, P_9, Q_9)\rho_9(t_9, a_9)da_9$$

Where $D_9$ is the mortality function and $n_9$ is the fertility function.

Put $_9$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_9}{dt}\right]_9 = x_9\left[\frac{dQ_9}{dt}\right]_9 - x_9\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_9}{dt}\right]_9 = x_3\left[\frac{dQ_3}{dt}\right]_9 - x_4\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_9}{dt}\right]_9 - x_6\left[\frac{dQ_3}{dt}\right]_9$$

Where $x_9, x_9, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_9$ and $g_9(P_9)Q_9$

Now when you add the term $(b_9Q_9)$ with $g_9(P_9)(Q_9)$ we are infect see is that $(b_9Q_9)$ is credited incrementally part of $g_9(P_9)(Q_9)$ or instant nearly, when subtracted from $\frac{dQ_9}{dt}$ , it means that $g_9(P_9)(Q_9)$ is slowly eaten by $\frac{dQ_9}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_9 = \mu_9(a_9) + (f_9)_9(P_9)_9 + (f_9)_9(Q)_9 \quad (2)$$

$$n_9 = \beta_9(a_9)h_9(Q_9), \quad (3)$$

Where $\mu_9$the age- dependent debit rate whit is $\mu_9(a_9)$ $\forall a > 0$; $(f_9)_9$ is the density- dependent death rate with $f_9(0) = 0$, $(f_9^9)_9(P_9)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_9)_9(P_9) = (k_9)_9$ ; $f_9$ is the functional response of creditors to prey with $[f_9]_9(0) = 0, (f_9^9)_9(Q) > 0, \forall Q$ , and $\lim_{Q\to\infty}(f_9)_9(Q_9) = [k_9]$ ; $\beta$ is the age dependent birth rate with $\beta(a)$ $, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^9(Q)$ $, \forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ .

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_9 + (\rho_a)_9 = -\mu_9(a_9) + (f_9)_9P_9) + \big(f_9)_9(Q_9)\big)(\rho_9)_9 \quad (4)$$

$$\rho_9(t_9, 0) = h_9(Q_9)\int_0^\infty \beta_9(a_9)\rho_9(t, a_9)da, \quad (5)$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(P_9)Q_9 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_9$ and $Q_9$ of (4-6) satisfy

$$0 \le \rho_9(t, a_9) \le \rho_9^*(t, a_9),$$

$$0 \le Q_9(t) \le Q_9^*(t),$$

Where $\rho_9^*(t, a_9)$ and $Q_9^*(t)$, are solutions of

$$(\rho_t^*)_9 + (\rho_a^*)_9 = -(\mu_9(a_9) + (f_9)_9(P_9^*) + (f_9)_9 Q_9^*))\, \rho_9^*, \quad (7)$$

$$\rho_9^*(t, 0) = \int_0^\infty \beta_9(a_9)\, \rho_9^*(t, a_9) da_9 \quad (8)$$

$$\frac{dQ_9^*}{dt} = -b_9 Q_9^* + g_9(P_9^*)Q_9^*, \quad (9)$$

With $P_9^*(t) = \int_0^\infty \rho_9^*(t, a_9) da_9$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_9^*(t)}{dt} = P_9^*(t)\big(r_9 - (f_9)(P_9^*) - (f_9)_9(Q^*)\big), \quad (10)$$

$$\frac{dQ_9^*(t)}{dt} = Q_9^*(t)(-b_9 + g_9(P_9^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_9(a) e^{-\int_0^a \mu_9(\tau)dt} e^{-ra}\, da_9 = 1. \qquad (12)$$

If $r < 0$, the trivial solution $P_9^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(t)dt}\, da_9 > 1, \qquad (H1)$$

Which guarantees $r_9 > 0$.

*Remark.* If $\beta_9(a) = [\beta_0]_9 a_9 e^{-\alpha_9 a_9}, \mu_9(a_9) \equiv \mu_9$ then the solution of (12) is

$$r_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_9]_9(P_9) = 0, [f_9](Q_9) = \lambda_9 Q_9\ h_9(Q_9) = 1, \qquad and\ \ g_9(P_9) = c_9 P_9$$

$$[\rho_9]_9 + [\rho_a]_9 = -(\mu_9 + \lambda_9 Q_9)\rho_9$$

$$\rho_9(t, 0) = \int_0^\infty [\beta_0]_9 a_9 e^{-\alpha_9 a_9} \rho_9(t, a_9)\, da_9$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + c_9 P_9 Q_9$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_9}{dt} = \omega_9 P_9 - \lambda_9 P_9 Q_9$$

$$\frac{dQ_9}{dt} = b_9 Q_9 - c_9 P_9 Q_9$$

Where $\omega_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$

THE EQUILIBRIUM $(\rho_0)_9(a_9), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_9(a_9), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_9 + (\rho_a)_9 = -\big(\mu_9(a_9) + (f_9)_9(P_9)\big)\rho,$$

$$\rho_9(t,0) = \int_0^\infty \beta_9(a_9)\rho_9(t,a_9)da_9 \qquad (13)$$

The solution of (13) has the form of

$$\rho_9(t,a_9) = Ce_9^{r(t-a_9)} - \int_0^a (\tau)d\tau - \int_0^9 (f_9)_9\big(P_9(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_9(t) = C_9 e^{r_9 t - \int_0^t (f_9)_9(P_9(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_9^{-9}(r)$ if $k_9 > r$. This is equivalent to

$$\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(\tau)d\tau} e^{(k_9)_9 a_9} da_9 < 1.$$

In fact, if we define

$$R_9(s_9) = \int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau)d\tau} e^{-s_9 a_9}\, da_9$$

$R(s)$ is a decreasing function. $R_9(r_9) = 1$ And $R_9(k_9)_9 < 1$ imply $r_9 < (k_9)_9$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_9(a_9), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(\tau)d\tau} e^{k_9 a_9} da_9$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_9 \equiv \int_0^\infty (\rho_0)_{99}(a_9)da_9$ *satisfies*

$$\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(\tau)d\tau} e^{-(f_9)_9 (P_0)a_9} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_9(a_9), 0)$

Let $u_9(t, a_9) = \rho(t, a_9) - \rho_0(a_9)$ Then

$$N_9(t) \equiv \int_0^\infty u_9\,(t, a_9)da_9 = \int_0^\infty \rho_9(t, a_9)da_9 - \int_0^\infty (\rho_0)_9(a_9)da = P_9(t) - P_0)_9$$

And

$$(u_9)_9 + (u_a)_9 = (\rho_t)_9 + (\rho_a)_9 - (\rho_0^9)_9(a_9)$$

$$= -\mu_9(a_9) + (f_9)_9(P_9) + (f_9)_9(Q_9))\,(u_9 + (\rho_0)_9) + \mu_9(a_9)(\rho_0)_9 + (f_9)_9(P_0)_9(\rho_0)_9$$

$$= -\mu_9(a_9)u_9 - (f_9)_9(P_9) + (f_9)_9(Q_9))\,u_9 - (\rho_0)_9\,(f_9)_9(P_9) - (f_9)_9(P_0)_9$$

$$+ (f_9)_9(Q_9)), \qquad (15)$$

Where $(\rho_0^9)_9(a_9)$ denotes $d(\rho_0)_9(a_9)/da_9$

Expansion $(f_9)_9(P_9)$ and $(f_9)_9(Q_9)$ at $(P_0)_9$ and 0 yield respectively

$$(f_9)_9(P_9) = (f_9)_9((P_0)_9 + N_9)$$

$$= (f_9)_9(P_0)_9 + (f_9^9)_9(P_0)_9N_9 + o_9(N_9), \qquad (16)$$

$$(f_9)_9(Q_9)$$

$$= (f_9^9)_9(0)Q + o_9(Q_9) \qquad (17)$$

Write $f_9(P_0) = f_{90}, f_9'(P_0) = f_{90}'$, and $f_9'(0) = f_{90}'$ and substitute (16) and (17) into (15) to obtain

$$(u_9)_9 + (u_a)_9 \approx -(\mu_9(a_9) + (f_{90})_9)u - (f_{90}^9)_9(\rho_0)_9(N_9) - (f_{90}^9)_9(\rho_0)_9(Q_9)$$

$$u_9(t,0) = \rho_9(t,0)(\rho_0)_9(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t,a) + (\rho_0)_9(a_9)da_9 - (\rho_0)_9\,(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t,a_9)da_9 + (h_9(Q_9) - 1)(\rho_0)_9\,(0) \qquad (18)$$

Expanding

$$h_9(Q_9) = 1 + (h^9)_9(0)Q_9 + 0(Q_9) \equiv 1 + h_0'Q_9 + 0Q_9$$

And substituting it into (18), we arrive at

$$u_9(t,0) \approx \int_0^\infty \beta_9\,(a_9)u_9(t,a_9)da_9 + (h_{a_9}')(\rho_0)_9Q_9$$

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PREDATOR-PREY POPULATION MODELS

The linearization is the following :

$$[u_9] + [u_9] = -(\mu_9(a_9) + [t_{90}]_9 u_9 - [\alpha_9]_9(a_9)N_9 - [\alpha_9]_9(a_9)Q_9$$

$$u_9(t,0) = \int_0^\infty \beta_9(a_9)u_9(t,a_9)da_9 + [\xi_9]Q_9 \qquad (19)$$

$$\frac{dQ_9}{dt} = -b_9 Q_9$$

where

$$[\alpha_9]_9(a_9) = [t_9^9 a]_9[p,o]_9(a_9)d2(a_9) = [f_{90}^9]_9(\rho\, o\,)_9(a_9)\xi = [\mathrm{h\ o}^9]_9[\rho\mathrm{o}]_9(\mathrm{o}) \text{ and } \mathrm{b}_9 = \mathrm{b}_9 - g_9(P_9)$$

If $g_9(P_0) > b$, clearly, $(\rho_0)_9(a_9), 0)$ is unstable and the predator population is persistent. Assume $g_9(P_0)_9 < b$.

Then, $Q_9(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_9 + u_a = -(\mu_9(a_9) + (f_{90})_9)u_9 - (\alpha_9)_9(a_9)N_9$$

$$= -(\mu_9(a_9) + f_{90})u_9 - \alpha_9(a_9)\int_0^\infty u_9(t,a_9)da_9$$

$$u_9(t,0) = \int_0^\infty \beta_9\,(a_9)u_9(t,a_9)da_9$$

Using the method developed by Cushing [2, 3], we suppose

$$u_9(t, a_9) = \tilde{u}_9(a_9)^{e^{\gamma(t-a_9)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_9}\, da_9$. Then,

A solution of the first differential equation for $\tilde{u}(a_9)$ gives,

$$\tilde{u}_9(a_9) = e^{-\int_0^a \mu_9(\tau_9)d\tau - (f_{90})_9 a}\left(\tilde{u}_9(0) - w\int_0^a e^{\int_0^\tau \mu_9(s_9)ds + (f_{90})}\tau_9\,(\alpha_9)_9(\tau_9)e^{\gamma_9\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9 - (f_{90})_9 a_9 - \int_0^a \mu_9(\tau_9)d\tau}\right)\tilde{u}_9(0)$$

$$+w\int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9 - (f_{90})_9 a_9 - \int_0^a \mu_9(\tau_9)d\tau}\int_0^a (\alpha_9)_9(\tau_9)e^{-\gamma_9\tau_9}e^{(f_{90})_9\tau_9 + \int_0^\tau \mu_9(s_9)ds}d\tau_9\, da_9$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^\infty e^{-\gamma_9 a_9 - (f_{90})_9 a_9 - \int_0^a \mu_9(\tau_9)d\tau}da\right)(\tilde{u}_9(0))_9$$

$$+\left(-1 - \int_0^\infty e^{-\gamma_9 a_9 - (f_{90})_9 a_9 - \int_0^a \mu_9(\tau_9)d\tau}\int_0^a e^{(f_{90})_9\tau_9 + \int_0^\tau \mu_9(s_9)ds + y\tau}\alpha_9(\tau)d\tau\, da\right)w = 0. \quad (23)$$

Define $Y(\gamma, a) = e^{-\gamma_9 a_9 - (f_{90})_9 a_9 - \int_0^a \mu_9(t_9)d\tau}$. Then

$$\left(1 - \int_0^\infty \beta_9\,(a_9) Y_9(\gamma_9, a_9) da_9\right) u_9(0)$$

$$+w \int_0^\infty \beta_9(a_9) Y_9(\gamma_9, a_9) \int_0^a (\alpha_9)_9(\tau)_9 (Y^{-9})_9(\gamma_9, \tau_9) d\tau_9\, da_9 = 0, \quad (24)$$

$$\left(\int_0^\infty Y_9(\gamma_9, a_9) da_9\right) \tilde{u}_9(0)$$

$$+\left(-1 - \int_0^\infty Y_9(\gamma_9, a_9) \int_0^a (\alpha_9)_9(\tau)_9 (Y^{-9})_9(\gamma_9, \tau_9) d\tau_9\, da_9\right) w = 0.$$

Equation (20) has a solution of the form of $u_9(t, a_9) = u_9(a_9) e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_9(0) \not\equiv 0$ and $\gamma, w$ in the complwx plane.

Let $I_9(\gamma_9, a_9) = \int_0^a (\rho_0)_9(\tau_9) Y^{-9}\,(\gamma_9, a_9) d\tau_9$. We obtain the characteristic equation

$$f_{90_9}^l \left(\int_0^\infty Y_9(\gamma_9, a_9) da_9\right)\left(\int_0^\infty \beta_9(a_9) Y_9(\gamma_9, a_9) I_9(\gamma_9, a_9) da_9\right)$$

$$+\left(1 - \int_0^\infty \beta_9(a_9) Y_9(\gamma_9, a_9) da\right)\left(1 + (f_{90}^l)_9 \int_0^\infty Y_9(\gamma_9, a_9) I_9(\gamma_9, a_9) da_9\right)$$

$$= 0. \quad (25)$$

Since

$$(\rho_0)_9(a_9) = (\rho_0)_9(0)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - (f_{90})_9 a_9}$$

$$I_9(\gamma_9, a_9) = \int_0^a (\rho_0)_9(\tau_9)e^{\gamma_9\tau_9 + (f_{90})_9\tau + \int_0^\tau \mu_9(s_9)ds_9}d\tau$$

$$= \frac{(\rho_0)_9(0)}{\gamma_9}(e^{\gamma_9 a_9} - 1)$$

Let

$$E_9(a_9) = e^{-(f_{90})_9 a_9 - \int_0^a \mu_9(\tau_9)d\tau_9}$$

Then $[Y_9]_9[\gamma_9]_9, a_9) = E_9(a_9)e^{-\gamma_9 a_9}$. by a simple calculation,

$$\int_0^\infty \beta_9(a_9)Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma}\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}\ da\right);$$

$$\int_0^\infty Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma_9}\left(\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}\ da_9\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_9(0)}{\gamma_9}[f_{90}^9]\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}da_9\ (\int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9)$$

$$+\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right).$$

$$\left(1 + \frac{(\rho_0)_9(0)}{\gamma_9}(f_{90}^{'})_9\left(\int_0^\infty E_9(\ a_9)da_9 - \int_0^\infty E_9(\ a_9)e^{-\gamma_9 a_9}da_9\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right)(1+\frac{(f_{90}^{'})_9}{\gamma_9}(\rho_0)_9(0)\int_0^{\infty}E_9(a_9)da_9)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (94), to

$$1-\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)da_9=1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\neq 0,$$

Since $\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)\int_0^{a}\alpha_9(\tau_9)Y^{-9}(\gamma_9,\tau_9)d\tau_9 da_9>0$. Therefore, the characteristic solution is

$$\gamma_9=-(f_{90}^{'})_9(\rho_0)_9\int_0^{\infty}E_9(a_9)da_9.$$

Since $f_9'(P_0)_9>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_9),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_9),0)$ which is related to the extinction of predators (debit summation types: that which acts as a resource)and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_9 + (\mu_9(a_9) + [f_9]_9(P_9^*) + [f_9]_9(Q_9^*))P_9^* = 0, \quad (26)$$

$$p_9^*(0) = h_9(Q_9^*)\int_0^\infty \beta_9(a_9)p_9^*(a_J)p_9^*(a_9)da_9 \quad (27)$$

$$g_9(P_9^*) - b_9 = 0, \quad (28)$$

Where $P_9^* = \int_0^\infty p_9^*(a_9)\, da_9$

Equation (28) yields $P^* = g^{-9}(b)$,

From (26),

$$p_9^*(a) = p_9^*(0)e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}$$

Substituting into (27), we need to solve

$$\left(1 - h_9(Q_9^*)\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}da_9\right)p_9^*(0) = 0$$

For $Q_9^*$

Fix $p_9^*(0)$ and $P_9^*$ Let

$$H_9(z_9) = 1 - h_9(z_9)\int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9}da_{,9}$$

And assume that

$$[k_3]_9 \int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(g_9^{-9}(b_9))a_9 - [k_9]_9 a_9}\, da_9 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty} [H_9(z_9)] > 0.$$

Clearly,

$$\frac{dH_9}{dz_9} = h_9(z)(\,[f_9^{'}]_9(z_9)\,) \int_0^{\infty} a_9\beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - f_9(P_9^*)a_9 - f_9(z_9)a_9} da_9$$

$-h^{'}(z_9) \int_0^{\infty} \beta_9(a_9)\, e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9} \quad da_9$

y the assumptions on $\left[f_9^9\right]_9 and\ [h^9]^9 dH_9/dz > 0$. Hence, if

$$H_9(0) = 1 - \int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9} da_9 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_9(0) < 0$ follows if $g_9^-(b_9) < [f_9^{-9}]_9(r_9)$, where $r_9$ is the solution of (12).

Indeed, since

$$R_9(s_9) = \int_0^{\infty} \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 e^{-y\frac{a}{9}} da_9}$$

Is a decreasing function of $s_9$ when $P_9^* = [g^{-9}]_9(b_9) < [f_9^{-9}]_9(r_9)$, that is $[f_9]_9(P_9^*) < r, R\big(f_9(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_9(g^{-9}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9\left(g_9^{-9}(b)\right)a_9} > 1. \qquad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_9^*(a_9), Q_9^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_9^*(a_9), Q_9^*)$ is more complicated. We again consider perturbations $m_9(t_9 a_9) = p_9(t_9, a_9) - p_9^*(a_9)$ and $W_9(t) = Q_9(t) - Q^*$.Then

$$M_9(t) \equiv \int_0^\infty m_9(t, a_9)da_9 = P_9(t) - P_9^*,$$

$$[m_9]_9 + [m_a]_9 = [p_9]_9 + [p_a]_9 - \left[p^*\right]_9(a_9) \approx -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9$$

$$-[f_{9*}^9]_9 p_9^*(a_9)M - [f_{9*}^9]_9 p_9^*(\mathrm{a})W_9,$$

$$m_9(t,0)p_9(t,0) - p_9^*(0) \approx h_* \int_0^\infty \beta_9(a_9)m_9(t,a_9)\,da + \frac{[h_*^9]_9}{[h_*]_9}\, p_9^*(0)W_9$$

$$\frac{dW_9}{dt} = \frac{dQ_9}{dt} \approx [g_*^9]_9 Q_9^* M_{,9}$$

Where $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $\left[f_{9*}^{'}\right]_9$ denoting $\left[f^{'}{}_9(P^*)\right]_9\left[f_{9*}^{'}\right]_9$ denoting $[f_9]_9$ $(Q_9^*,$ $h_9^*$ denoting $h_*^{'}$ denoting $h_9^{'}$ $(Q_9^*)$, and $\left[g_*^{'}\right]_9$

Denoting $[g^{'}]_9(P_9^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_9]_9 + [m_a]_9 = -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9 - [f^{'}_{9*}]_9[p_9^*](a_9)M_9 - [f^{'}_{9*}]_9 p_9^*(a_9)W_9$$

$$m_9(t,0) = [h_*]_9 \int_0^\infty \beta_9(a_9)m_9(t,a_9)da_9 + \frac{[h_*^9]_9}{[h^*]_9}[p^*]_9(0)\, W_{,9} \qquad (30)$$

$$\frac{dW_9}{dt} = [g_*^9]_9[Q_9^*]M_9$$

Integrating $[m_9]_9(t,a_9)$ with respect to a from 0 to $\infty$, we have

$$\frac{dM_9(t)}{dt} = \int_0^\infty [mt]_9(t,a_9)da_9$$

$$= -\int_0^\infty [m_a]_9\,(t,a_9)da_9 - \int_0^\infty \mu(a_9)m(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + [f^{'}_{9*}]_9 \int_0^\infty p_9^*(a_9)da_9\right)M_9 - [f_{9*}^9]_9 \int_0^\infty p_9^*(a_9)da_9 W_9$$

$$= m_9(t,0) - \int_0^\infty \mu_9(a_9)\, m_9(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + [f^{'}_{9*}]_9 P_9^*\right)M_9 - [f^{'}_{9*}]_9 P_9^* W_{.9}$$

From (30), we arrive at the equations

$$\frac{dW_9(t)}{dt} = \int_0^\infty ([h_*]_9\beta_9(a_9) - \mu_9(a_9))m_9(t,a_9)da_9$$

$$-\left([f_{9*}] + [f_{9*}]_9 + [f_{9*}']_9 P_9^*\right) M_9 + \left(\frac{[h_*']_9}{[h_*]^9}\, P_9^*(0) - f_{9*}' P_9^*\right) W_9,$$

(31)

$$\frac{dW_9(t)}{dt} [g'_*]_9\, Q_9^*\, M_{.9}$$

By the mean value theorem,

$$\int_0^\infty ([h_{"*}]_9 \beta_9(a_9) - \mu_9(a_9)) m_9(t, a_9) da_9 = ([h_*]_9 \beta_9(\theta_9) - \mu_9(\theta_9)) M_{,9}$$

Where $0 < \theta_9 = \theta_9(t) < \infty$.

And if a $^9_9$ [a $^9_9$]2<4 [a $_9$]2<4 a $_9$e,

Re $\lambda_9 [\lambda_9]_9 = \text{Re}\lambda_9 [\lambda_9]_9 = -\frac{9}{9} a_9 [a2\ ]_9 < 0.$

Let ⊿=max$_{9<u<\infty}$ |h $_*$ [h $_*$] $_9$β $_9$($a_9$) − $\mu_9(a_9)$|. Then|$n_9$(t)| ⊿∀t . Hence we obtain the following result.

**THEOREM 4**

Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium(ρ $^*$ (a), $\boldsymbol{Q}^*$) is asymptotically stable if either a $^9_9$ <4 a $_9$ e and [a $_9$] > 1⊿,or [a $^9_9$] 4[a $_9$]e and (a $_9$ − $\sqrt{(a\ ^9_9)_9 - 4(a_9)_9 e}$ > 1⊿

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+9}\right)$ and $\beta_9(a_9) = ae^{-ma}$ Then $\Delta=[\mu_0] + 2$ Simple calculation leads to

$$\int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9) + d\tau_9 da} = \frac{1}{(m + \mu_0 + 1)^9} + \frac{1}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_9 + [\mu_0]$ 0.5. Assume$[f_9](x_9)= K_9 x_9$ Then the hypothesis (H9) is satisfied We take the conversion rate function $g_9(x_9)= [0x_9]/(2 + x_9)$; the functional responses $[[f_9](x) = [0]\, x_9/(1 + x_9)$and $h_9(x_9)=(1+\omega_9 x_9)/(1 + x_9)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_9 + [\mu_0] + [K][P_0])}{[m_9]_9 + (\mu_0)_9 + [K][P_0])_9^3\ [K]_9[P_0]_9^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_9$i e $[P_0] < b/(\delta_9 - b_9)$then the equilibrium$(P_0)$ $[P_0]_9(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_9$

$[f_{9*}] = K_9\ P^*\ [f_{9*}] = \frac{\sigma Q_9^*}{9+Q_9^*}$ $[h_*] = \frac{9+\omega_9 Q_9^*}{9+Q_9^*}$

$[f_{9*}] = K_9\ [f_{9*}^9] = \frac{\sigma_9}{(9+Q_9^*)^9}[\ g_*^9] = \frac{\sigma_9}{(9+P_9^*)^9}$ $[h_*^9] = \frac{\omega_9 - 9}{(9+Q_9^*)^9}$

In order to have (H3) and (H4) hold,$m, \mu_0$, K, b, $\delta$, and $\omega$ need to satisfy

$$\frac{\omega_9}{(m + [\mu_0]_9 + 1 + \frac{K_9 b_9}{\delta_9 + b_9} + \sigma)^9} + \frac{1\omega_9}{([m_9]_9 + [\mu_0]_9 + 1 + \frac{K_9 b_9}{\delta_9 + b_9} + \sigma_9)^3} < 1$$

And

$$\frac{1}{(m_9 + [\mu_0]_9 + 1 + \frac{K_9 b_9}{\delta_9 - b_9})^9} + \frac{1\omega_9}{(m_9 + [\mu_0]_9 + 1 + \frac{K_9 b_9}{\delta_9 - b_9})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_9 + [\mu_0]_9 + 1 + \frac{K_9 b_9}{\delta_9 - b_9} + (m_9 + [\mu_0]_9 + \frac{K_9 b_9}{\delta_9 - b_9} + \sigma_9)\, Q_9^*)^3$$

$$= m_9 + [\mu_0]_9 + 3 + \frac{K_9 b_9}{\delta_9 - b_9}(1 + Q^*)(1 + \omega_9 Q_9^*$$

Now

$$\begin{bmatrix} a & 1 \\ 1 & \end{bmatrix} - 4\begin{bmatrix} a & \\ 9 & \end{bmatrix} e = \left(\frac{1K_9 b_9}{\delta_9 - b_9}\right)^9 + \frac{4K_9 b_9 \sigma_9 Q_9^*}{(\delta_9 - b_9)(9 + Q_9^*)} + \frac{\sigma_9^9 Q_9^{*9}}{(1 + Q_9^*)^9}$$

$$\frac{4\, b_9 \sigma_9 (\delta_9 - b) Q^*}{\delta_9 (1 + Q_9^*)^9} - \frac{4\, (\delta_9 - b_9)^9 (1 - \omega_9) Q_9^*}{\delta_9 (1 + \omega_9 Q_9^*)(1 + Q_9^*)} P * (0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_9\, ae^{mu}_9\ \ [\mu_0]_9 - \frac{a_9}{a_9 + 9}$$

Hence if $[a_9^9 > 4(a_9)_9 e$ and

$$[a_9] = \frac{1K_9 b_9}{\delta_9 - b_9} + \frac{\sigma_9 \, Q_9^*}{1 + Q_9^*} > 1\Delta$$

Some of the examples to which the model has been applied are:

**New look at the galaxy-shaping power of black holes**

(1) "It creates (eb) these bubbles, and the bubbles keep the plasma hot.
(2) That's what prevents galaxies from (e) becoming even bigger than they are now."
(3) Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.
(4) In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.
(5) Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.
(6) "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.
(7) "It's like nature's thermostat, which keeps these galaxies from (e) growing.
(8) If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."

**A new look at the galaxy-shaping power of black holes**

(9) An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the
(10) McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.
(11) Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.

**Doomed satellite missions**

(12) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.
(13) At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.

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| (14) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up. |
| (15) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies. |
| (16) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said. |
| (17) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people |
| (18) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole." |
| (19) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma. |
| (20) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website |

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

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| (21) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming. |
| (22) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes. |
| (23) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking. |
| (24) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses. |
| (25) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well. |
| (26) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma. |
| (27) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6. |
| (28) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of $164 \pm 10$ kilometres per second in the region 30–60 kiloparsecs from the central nucleus. |

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| (29) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (30) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (31) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (32) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (33) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (34) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (35) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (36) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (37) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (38) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (39) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (40) We know dark matter exists only from (e) its gravitational effects. |
| (41) It holds together galaxies and (e&eb) galaxy clusters. |
| (42) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (43) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |
| (44) They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (45) X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (46) They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (47) This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (48) An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |

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| (49) Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (50) "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (51) The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (52) Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (53) Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (54) "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (55) Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (56) "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (57) Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories. |
| (58) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen. |
| (59) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer. |

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| (60) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants. |
| (61) These were treated as (=) "aware" and "unaware" trials, respectively. |
| (62) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it. |
| (63) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another. |
| (64) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion. |
| (65) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness. |
| (66) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said. |
| (67) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness." |
| (68) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner. |
| (69) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain. |
| (70) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said. |
| (71) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

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| (72) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter. |
| (73) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined. |
| (74) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions. |

(75) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.

(76) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.

(77) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.

(78) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.

(79) Noda and his colleagues instead studied the change in emissions over (e&eb) time.

(80) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.

(81) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).

(82) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.

(83) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.

(84) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(85) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking "palm tree" organization.

(86) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society

(87) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible?

(88) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem.

(89) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other.

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| (90) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible. |
| (91) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges. |
| (92) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century. |
| (93) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching. |
| (94) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom. |
| (95) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth. |

**Spheres on trees**

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| (96) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration. |
| (97) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement. |
| (98) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum. |
| (99) By visualizing the problem in this way, the researchers noticed that these particular graphs don't (e) have lots of branches extending from lots of (e) other branches. |
| (100) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure. |
| (101) The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network. |
| (102) As a result, it provides (eb) an important clue for finding the global minimum. |
| (103) It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum. |
| (104) This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large. |

**Thompson's problem 2**

(105) By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society

(106) The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima.

**(107)** The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node.

**Implications for Thomson's and Smale's problems**

(108) Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before.

(109) "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around."

(110) As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem.

(111) "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view."

(112) In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges.

(113) From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear.

(114) "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK.

(115) "For larger numbers, there will likely be (=) multiple funnels.

(116) Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum."

(117) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.

(118) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.

(119) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org

The nontrivial equilibrium is stable.

IF we define $[a_9]_9 = [f_{9*}]_9 + [f_{9*}]_9 + [f_{9*}]\, P^*_9$

$$[a_9]_9 = \left[f^{9}_{9*}\right] P^*_9 - \frac{h^9_*}{h_*} P^*(0)$$

$$e=[\, g^9_*]_9\, Q^*_9$$

$$\eta_9(t) = \left[h_*\right]_9 \beta_9 \left(\theta_9(t)\right) - \mu_9(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\left[a_9\right] M_9 - [a_9]_9 W_9 - [n(t)] M_9$$

$$w_9(t) = eM_9$$

Write

$$A=\begin{pmatrix} -[a_9]_9 - [a_9]_9 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_9 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_9 \\ W_9 \end{pmatrix}$$

The solution of (32)is

$$X_9(\mathrm{t}) = \mathrm{e}^{A,t}\,[\mathrm{X}]_9(0) + \int_0 \mathrm{e}^{At}$$

**Note: There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all**

**production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities**

**would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being**

**added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the**

**parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module**

**LEN NUNNEYS**

# Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

## Module numbered one

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay

and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + \ N\frac{dG}{dR} - \frac{dB}{dR}$$

$$> 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1965); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be referred to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 6985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_6}{dt} = [N(t)]_6[(F_6)(R(t)]_6 - [M]_6[R(t)]_6$$

$$\frac{d[R(t)]_6}{dt} = [B]_6\,[(R)_6(t-T_6)] - [D]_6[(R(t)]_6 - [N(t)]_6\left[G\left(R(t)\right)_6\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_6}{dt} = a_6[N(t)]_6 - a_6[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_6 = a_3[N(t)]_3 - a_4[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_6 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

irst, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.6b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 6a is at the

**TIME DELAY**

FIG. 6. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in

the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.6$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

## DELAY DUE TO PREDATOR MATURATION

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_6}{dt} = N_6(t - T_6)F_6(R_6(t - T_6) - N_6(t)M_6((R_6(t) \qquad (5)$$

$$\frac{dR_6(t)}{dt} = B_6\big(R_6(t) - D_6(R_6(t))\big) - N_6(t)G_6(R_6(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_6}{dR_6}+\frac{N_6 dG_6}{dR_6}-\frac{dB_6}{dR_6}>\frac{N_6 G_6}{2M_6}\left[\frac{dF_6}{dR_6}+\frac{dM_6}{dR_6}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_6/dR_6$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_6$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_6$) and the other derived from the non-delayed or "immediate" processes($I_6$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

ABSOLUTE STABILITY

$$L_6=\frac{N_6 G_6}{2M_6}\frac{d(6nF_6)}{dR_6}(\geq 0) \quad (6a)$$

$$I_6=-\frac{N_6 G_6}{2M_6}\frac{d(6nM_6}{dR_6}(\geq 0) \quad (6b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_6>L_6-I_6$, must be modified ahe $dM_6/dR_6$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_6>L_6-I_6, \quad \text{when} \quad P_6^2/4>I \quad (8a)$$

$$P_6>L_6-I_6+\frac{(I_6-P_6^2}{(L_6-I)} \quad \text{when} \quad P_6^2/4<I. \quad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

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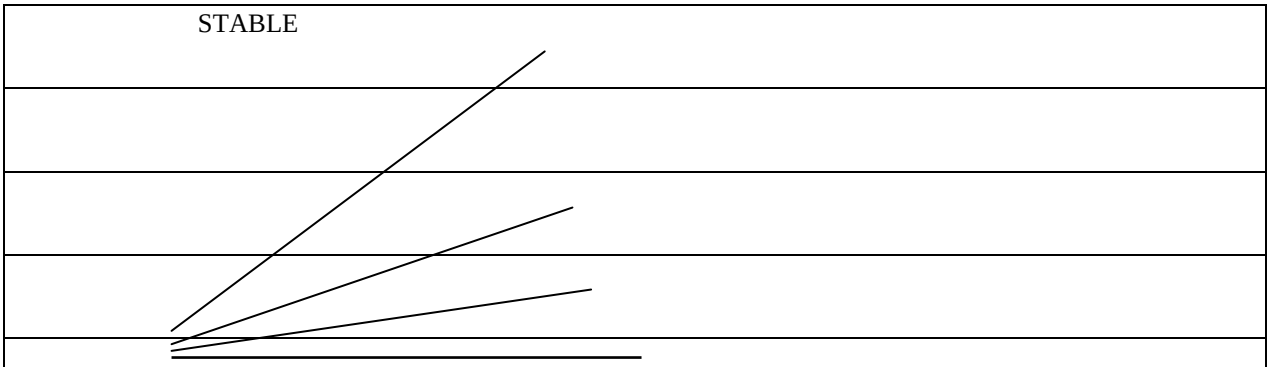



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

Can never completely offset the lagged effect,$L_6$; condition (8b) ensures that the total debit regulation, $P_6$, must always be positive even when $I_6$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_6$ and $L_6$. As $I_6$ increases relative to $L_6$ the zone of stability is enlarged; however,even when $I_6$is zero a significant zone of stability exists.

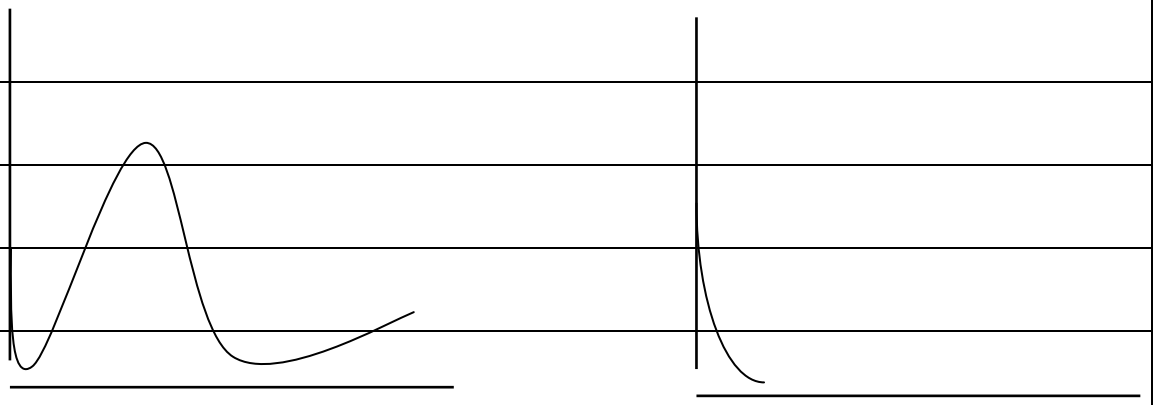
It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay

and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_6$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



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Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (6). Parameters in each case M=6; $dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6 = 0.05$ For (a) $NGdF/dR = 0.3, -N_6G_6dM_6/dR_6 = 0.6$; for (b)$NG_6dM_6/DR_6 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_6t}{dt} = N_6(t - T_6)F_6\big(R_6(t - T_6)\big) - N_6(t)M_6\big(R_6(t)\big)$$

(9)

$$\frac{dR_6t}{dt} = B_6\big(R_6(t-T_6)\big) - D_6\big(R_6(t)\big) - N_6(t)G_6\big(R_6(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) and the effective prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_6/dR_6$ is a measure of the control acting through the debit delay and setting $dB_6/dR_6 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_6/dR_6$ is a measure of the control acting through the credit delay although setting $dF_6/dR_6 = 0$ does not negate the effect of a predator delay; the value of $M_6$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the

delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_6/dR_6$and normalized by $M_6$ (see6a). The absolute stability of such a system is summarized in Fig. 6 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

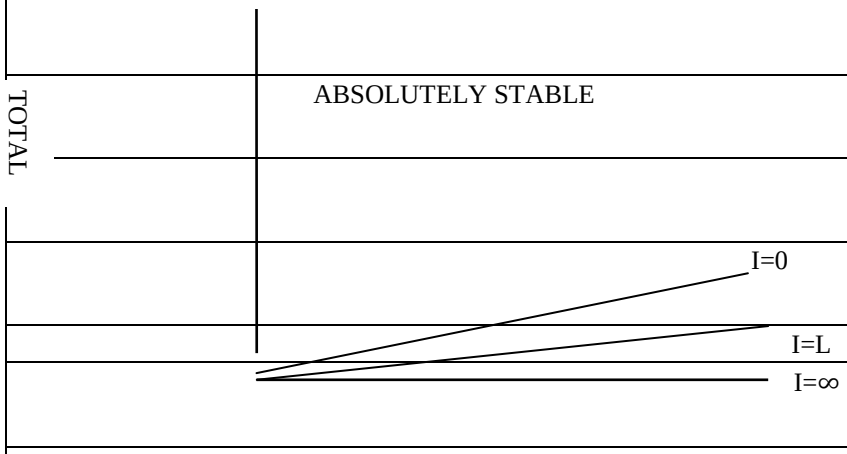



FIG6. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_6/dR_6 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.

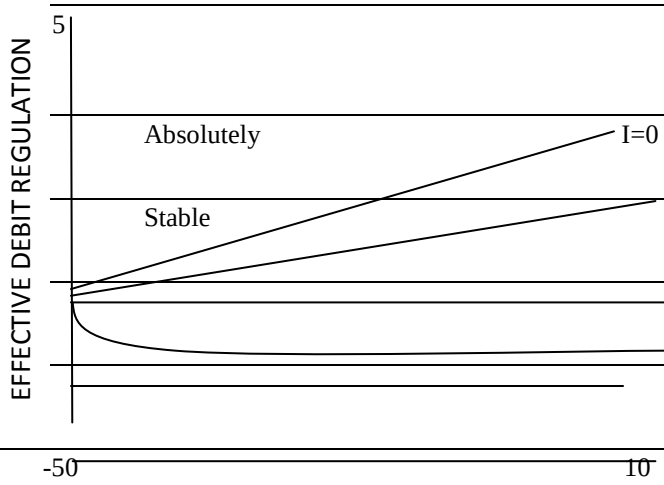



Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_6/dR_6 + N_6 dG_6/dR_6 + dB_6/dR_6)/M_6$). The immediate ($I_6$) and lagged ($L_6$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_6$(see 6b), is increased relative to the lagged part, $L_6$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_6 > I_6$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_6$is somewhat large than $I_6$ then the condition constraining prey regulation of recruitment ($dB_6/dR_6 < 0$) is

$$\frac{dD_6}{dR_6} + N_6\frac{dG_6}{dR_6} + \frac{dB_6}{dR_6} > \frac{(L_6 - I_6)M_6}{2} \qquad (10)$$

If $L_6 < I_6$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_6$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_6$

The - $D_6$ positions are defined by letting $\lambda_6 = iZ_6$ and separating (A1.2) into real and imaginary parts:

$$Z_6^2 + V_6 Z_6 \int_0^\infty \sin(Z_6 T_6) f(T_6) dT_6 - (X_6 + Y_6) = 0$$

$$Z_6\left(U_6 - V_6\int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6\right) = 0$$

Noting that $X_6 + Y_6 = 0$ is we assume X+Y=0.It follows that $Z_6 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_6 = V_6\int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6$$

If $U_6 > |V_6|$then given the constraints of $f(T_6)$, $(A1.5)$*can never be satisfied Co*nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A6.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CREDIT CONSOLIDATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_6^2 + (U_6 - V_6 + M_6)\lambda_6 - M_6\lambda_6 e^{-T_6\lambda_6} + [(U_6 - V_6)M_6 + Y_6] - [((U_6 -$$

$$V_6M_6-]e^{-T_6\lambda_6}=0$$

(For definitions see Appendix6) We follow the technique outlined in Appendix 6 and let $\lambda_6 = iZ_6$ and separate (A2.6) into real and imaginary parts

$$\boldsymbol{Z}_6^2 + \mathrm{M}_6\mathrm{Z}_6\sin(\mathrm{Z}_6T_6) - (A_6M_6 + Y_6) + (A_6M_6 + X_6)\cos(\mathrm{Z}_6T_6) = 0$$

$$(A_6M_6)\mathrm{Z}_6 - \mathrm{M}_6\mathrm{Z}_6\cos(\mathrm{Z}_6T_6) + (A_6M_6 - X_6)\sin(\mathrm{Z}_6T_6) = 0$$

Where $A_6 = \mathrm{U}_6 - \mathrm{V}_6$Using (A2.3), substitute for $\boldsymbol{Z}_6^2$ and $\mathrm{Z}_6$in (A2.2):

$W_6^2(X_6^2$-3 $M_6\ A_6\ X_6$- $M_6^2(X_6$- $Y_6) + 2\ M_6^3\ A_6$+2 $A_6^2\ M_6^2]$

$$+W\left[MA^3 - A^2\mathrm{X} + 2\mathrm{MA}(2\mathrm{X} + \mathrm{Y}) - 2X^2\right] + A^2(\mathrm{X} + \mathrm{Y}) = 0$$

*where* $W_6$=1-cos $(\mathrm{Z}_6T_6)$.

In the absence of a delay the system is stable for $A_6 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_6 > (\mathrm{X}_6 - Y_6)/2\ M_6$ Thus a necessary condition for absolute stability is

$A_6 > [\ \mathrm{X}_6 - Y_6/2\ M_6$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_6$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_6$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_6, A_6) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_6 \le 2$ *is* of interest First consider the extreme values

$$Q_6(0, A_6) = A_6^2 (X_6 - Y_6)$$

$$Q_6(2, A_6) = (A_6 + 2M_6)^2 (2A_6 M_6 - (X_6 - Y_6))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_6}{\partial A_6} = W_6^2 \left(2M_6^3 + 4A_6 M_6^2 - 3M_6 X_6\right) + W_6 \left(3M_6 A_6^2 - 2A_6 X_6 + 2M_6 \left(2X_6 + Y_6\right)\right)$$

$$+2A_6 \left(X_6 + Y_6\right)$$

$$=2M_6^2 W_6^2 (M_6 + A_6) + M_6 A_6^2 W_6 + (2A_6 M_6 - X_6 + Y_6)(A_6 W_6 + M_6 W_6^2$$

$$+ (2 - W_6)((2M_6 X_6 + M_6 Y_6) W_6 + A_6 (X_6 + Y_6)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_6$quadratic in $W_6$touches the $W$ –axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_6$has no minimum in the range $0 \le W_6 \le 2$.Defining

$Q_6(W_6, A_6) = a_6 W_6^2 + c_6$ *where* $a_6$ $b_6$ *and* $c_6$ Are functions of $A_6$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_6 - b_6 < 0$$

And

$$b_6^2 > 4a_6c_6$$

$$Case(i): Y_6 = 0$$

$Let\ Y_6 = 0$ And try the substitution A= $X_6 = /2M_6\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_6 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_6 X_6 /$ $2M_6$) is positive for *W<2. Since* $Q_6$ increases with $A_6$(see A2.9), increasing $A_6$will generate no further solutions of $Q_6$=0. Hence the condition A>X/2M defines absolute stability when $Y_6 = 0$.

$Case\ (ii): Y_6 = y_6\ X_6$ ,

$Let\ Y_6 = y_6\ X_6$And try the substitution $A_6 = X_6/(2 + A_6\ y_6)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_6/(2 + n_6 y_6)$ we have

$$W_6^2 y_6\left[X_6 n_6 \frac{(1 + n_6 y_6)}{(2 + n_6 y_6)} + M_6^2\left(2 - n_6 + n_6 y_6\right)\right]$$

$$-W_6 X_6 \left[\frac{X_6(1+n_6 y_6)}{M_6^2(2+n_6 y_6)^2}+2y_6\left(n_6-1\right)+\frac{X_6^2(6+y_6)}{M_6^2(2+n_6 y_6)}\right]=0.\ (A2.12)$$

Since $n_6$ and $A_6$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_6$

Defining $E_6=(2+n_6 y_6)$, *the condition*(A2.66?)*can be written as*

$[\frac{X_6(E_6-6)}{E_6^2}]^2[\frac{X_6^2}{M_6^2}-4y_6 E_6^2\frac{X_6}{M_6^2}+4y_6(n_6-2)E_6^3]$ ?0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_6/W_6^2)$ can be expressed as

$2E_6[y_6E_6+(y_6E_6(6+y_6)(E_6-n_6))^{6/2}]>\frac{X_6}{M_6^2}>2E_6[y_6E_6-y_6E_6(6+y_6)(E_6-n_6))^{6/2}]$

**ABSOLUTE STABILITY**

Hence if $X_6/M_6^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0?$W_6$ ?2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_6-6)}{E_6^2}\left(\frac{X_6}{M_6^2}\right)^2-2y_6\frac{(E_6+n_6^2y_6)}{E_6}\left(\frac{X_6}{M_6^2}\right)-4y_6(E_6-n_6)\ ?0$$

And has the solution

$$\frac{X_6}{M_6^2} > \frac{E_6}{(E_6-6)}[y_6(E_6 + n_6^2 y_6) + [y_6^2 (E_6 + n_6^2 y_6)^2 + 4 y_6 (E_6 - 6)(E_6 - n_6)]^{6/2}]$$

Given $E_6 > n_6$ this lower limit to $X_6 / M_6^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_6 / M_6^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_6$ in the appropriate range and this limit is further lowered by (A2.64)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_6 < n_6$ $(ie., 2A_6 M_6^2 < X_6 - Y_6)$ then (A2.6)and (A2.8)show that $Q_6$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_6}{M_6^2} > 2E_6[y_6 E_6 - (y_6 E_6 (6 + y_6)(E_6 - n_6)]^{6/2}]$$

This equations throws light on the individual general ledger and Nature's general ledger as we call it : see various papers on the issue. powers, which are the energies of Matrka we have already discussed, are not the only obstacles the yogi must overcome. He must, for example, also resist the temptation to rest content with the miraculous yogic powers (siddhi) he acquires in the course of his spiritual develop- ment. Again to do this he must practice Yoga. Similarly, in order to pervade the Fourth State gradually through the other states in the manner proper to practice at the Individual level, the method is the same. He must practise the higher yoga of the Tantras which, turning his mind inwards and freeing it from discursive representations, allows him to penetrate into the Supreme Principle. 312 Once the yogi has attained this contemplative state, his

main problem is to make it permanent. In the introverted state the gross external movement of the breath is suspended and with it the activity of the intellect, mind, individualised consciousness, powers of the senses and the ego. 313 When the yogi rises out of this state, he is liable to fall again into the lower order of creation generated by Maya if he does not maintain his awareness of the higher reality he has experienced and allows his awakened, illumined insight to be obscured by the dream-like vision of thought-constructs. 314 Naturally, the yogi must rise out of the introverted condition of suspension. It is inherent in the very nature of reality that it should move out of itself. 315 Pure, universal consciousness initially transforms itself into the vital breath 316 charged with the impression iyasana) of the power of awareness attained through introversion. By attending to the pulse (spanda) of the breath as it moves out of the absolute, the yogi can develop an intuitive sense of the inherent unity of all he will perceive in the mental and physical spheres created by the outpouring of consciousness. In this way he realises that his own nature is everywhere present in all he perceives and that all things thus reside within him. Blessed with this insight his consciousness remains free and unlimited even at the individual level where the breath, mind, senses and body are active. If the yogi fails to do this, he finds himself once again beset by the strictures of his embodied existence and must, as before, try to pervade all his other states of consciousness with the aesthetic delight (rasa) and wonder (camatkara) of the Fourth State he experienced in contemplation. Again this means that he must strengthen his pure, empowered awareness that his universal nature manifests as all things. 317 In this way he discovers Siva's presence in every sphere of individualised consciousness ranging from the breath to outer objectivity. The yogi's mind then becomes tranquil and undistracted because wherever it may wander, the yogi perceives only Siva, his authentic nature. 318 Consciousness is thus freed of all external referents and the yogi's subjectivity is purified of all identification with the body or anything else that belongs to the objective sphere. The yogi then becomes detached from the opposites of

pleasure and pain and is transcendentally free (kevalin). 319 yogis are again, however, liable to fall if he allows himself to get entangled in the play of opposites. This fall is more serious than the others because, although he is caught by the confining restrictions of individual- ised consciousness as before, he is now also affected by karma. Fleeing from pain in the pursuit of pleasure he is bound to act (karma) to minimise one and maximise the other and so is thrown down to the lowest level of embodied subjectivity (sakala). In order to regain his lost state, he must ascend gradually, by Siva's grace, from one order of subjectivity to the next and so free himself progressively of the limitations of the lower levels to gain the greater freedom and expansion of the higher. As he progresses, the objective sphere also evolves from the grossest perceptions of physical objects outside the lowest order to subjectivity, through to the subtler inner, mental perceptions to finally reach the order of subjectivity that contains objectivity within itself and is free to externalise it at will. 320 The degree to which this process develops depends, as before, on the yogi's awareness of the Fourth State. In consonance with the general principle that the remedy should suit the defect, the yogi is instructed to seek this higher state of consciousness in the wonder (camatkara) or delight (ananda) he feels in moments of intense physical pleasure. At first he experiences this subtle consciousness for an instant in the subjective sphere. If he manages to catch hold of it, it becomes more intense as the cognitive and objective spheres are also gradually pervaded and vitalised by it. Occasions for this practice are, for example, the sense of satisfaction one feels after a good meal or the aesthetic delight one experiences when listening to good music or the pleasure of sexual union with the Tantric consort or even solitary sexual excitation. In these moments of delight the yogi can penetrate momentarily into his own authentic Siva-nature (sambhavavesa) through the empowered contact (saktasparsa) m he makes with it in the freedom of the pure subjectivity of the Fourth State. 322 If the yogi develops his awareness of this higher level of consciousness and maintains it, he eventually experiences it constantly. 323 Clearly, what

prevents the yogi from attending to his state of consciousness rather than the circumstances which induce it is the craving for pleasure (abhila§a) born of ignorance — the source of every impurity which clouds consciousness. Craving directs the yogi's attention towards outer, worldly things and so he is caught in the net of thought- constructs. 324 To free himself of his worldly desires and reverse this binding extroversion, the yogi must eradicate its cause. To be freed of all the ups and downs of the path and no longer be tormented by the possibility of a fall, the yogi must see reality perfectly and completely. This insight is itself liberation and the moment it dawns the yogi is instantly freed. This sudden realisation is the goal of Tantric Yoga. E Tantra declares: "He who perceives reality directly, even for the brief moment it takes to blink, is liberated that very instant and never reborn again." 325 Although the yogi's body and mind continue to function as before, they are like mere outer coverings 326 which contain, but do not obscure, the mighty, universal consciousness which operates through them. The yogi's body is (=) the universe, the senses (=) the energies that (eb, eb+)vitalise it, his mind (=) Mantra, the rhythm of his breath (eb) the pulse of time and his inner nature pure, dynamic consciousness. Raised above all practice, and hence all possibility of falling to lower levels, the yogi realises that he has always been free 327 and that his journey through (e&eb) the **dark land of Maya was nothing but a dream, a construct of his own imagination** The channels of the senses by which the yogi enjoys the pleasure of the objects of sense are those through which he fills the three worlds with the consciousness emitted from his own Heart. 78 **The Doctrine of Vibration T**he senses are outer, subsidiary cycles of conscious energy (anucakra), revolving around the core or primary cycle (mukhyacakra) of pure consciousness. They scintillate with the diversity of sensations and resonate in the field of their awareness like the feathers of a peacock unfurled and folded back in the ecstasy of its dance. 79 Ultimately it is Siva Himself Who, for His own pleasure, 80 extends and withdraws them, 81 not the individual soul. The outward movement of the senses is the creative flow

of consciousness from the subject who, like an ocean, is the source of the waves of their activity and the final resting place of the flow of the rivers of sensations and thought-forms that unfold through the arising of the power of each of the senses. 82 The Stanzas on Vibration declare: Indeed the individual soul does not activate the impulse of the will [which directs the body's activity] by himself alone but through his contact with [his] own [inner] strength {bald) made in such a way that he identifies with it [thus acquiring its power]. 83 The close contact between the senses and universal consciousness vitalises them and sets them in motion. They share in the properties of consciousness, just as a heated iron ball can burn, cook or give out light. 84 Spanda doctrines does not agree with the Sarnkhya view that the individual soul (purusa) alone stimulates activity in the objective sphere (prakrti) by the contact he makes with it. The soul can do this only because universal consciousness animates him. Moreover, consciousness, which has the power to impel this activity, cannot itself be inactive as the Sarnkhya maintains. The influence it has on lifeless matter implies an active exertion. 85 This activity is the subtle, inner outpouring of the Self (svatmocchalatta) within the Self. Outside the confines of time and space, in the centre of the circle of the senses, this inner activity initiates the expansion (prasara) of the flow of conscious energy. As the pure exertion (udyoga, udyama) of Being (bhava), it is the source of the initial unfolding (unmesa) of Becoming. 86 Ksemaraja explains: [By] 'exertion' (udyama) is [meant] the emergence of the supreme [level of] intuition (pratibha) which is the sudden outpouring of the [immanent and transcendental] unfolding of the reflective awareness of pure consciousness. 87 The direct, intuition (pratibha) of reality, free of thought-constructs, is the supreme power (parasakti) of consciousness. The yogi can gain Divine **Body and Sacred Circle of the Senses 153** this pure intuitive awareness by grasping the initial expansion (unmesa) of consciousness which engenders the emergence of thought and the activity of the senses. It can be discovered in the subtle Centre, between one thought and the next. The Stanzas on Vibration teach: The

expansion of consciousness that takes place when one is engaged in a single thought should be known to be the source from whence another arises. One should experience that for oneself! 88 This does not mean that the unfolding expansion of consciousness {unmesa) is an intermittent, transitory occurrence. 89 As the source of all thought, including the notions of past, present and future, 90 it stands outside the confines of time. Consciousness is in a state of perpetual expansion. It is an ever developing process, a tiny part of which we observe as the flow of events around us and the thoughts and feelings within us. The totality of this process can never be grasped objectively. Objectivity is a state of fragmentation (bheda) in which discrete elements are discernible and mutually distinguished as individual, specifically definable entities. It is the realm within which thought and language operate. To grasp reality in its completeness, we must go beyond the partial representations of thought and speech. To experience the primordial source and basis of all things, we must pierce through the outer periphery of thought and plunge into the Centre, to discover the instant in which thought, and with it, the sense of diversity, initially emerges. Ksemaraja cites a practice from the Vijnanabhairava through which the yogi can achieve this state of awareness: Checking [the movement of] attention (cit) once it has quit its object, it must not [be allowed] to move to any other; contemplation (bhavana) then blossoms forth by [experiencing] the state in the centre between them (tanmadhyabhava). 91 Left unattended and undisciplined, awareness spontaneously shifts from one object of thought to (e&eb) the next. The transition from one to the other entails a movement through (e&eb) a state of pure indeterminate (nirvikalpa) awareness divested of all thought-forms and perceptions. The yogi who succeeds in checking the movement of his attention, experiences within himself the subsidence of thought into the emptiness of the Centre which unfolds full of the powerful pulsation of consciousness within which all objectivity merges and becomes one 92 in the trans- cendental outpouring (lokottarollasa) of its aesthetic rapture. 93 turning in on himself, the yogi must exert his power of awareness to (e) discern

the vitality of (e) his own blissful Spanda nature. Just as a distant object becomes (=) more clearly visible by (e) exerting an effort to (e) perceive it, in the same way, through (e&eb) practice, the yogi's own Spanda nature becomes (=) progressively more evident to him, 94 The effort the yogi exerts to free himself of ignorance is essentially one with the force of awareness that impels the senses and mind. This inner exertion is the basis of every means to liberation. Ultimately bestowing the bliss of self- realisation, it carries the yogi along the path to salvation with all its twists and turns like wind which blows a cloth in various directions. 95 As the basis and ultimate goal of all yogic practice, this exertion is the pure wonder (camatkdra) inspired by the realisation of the essential identity of consciousness (cit) and being (sat). 96 As such it is the divine radiance (sphuratta) of the Heart of one's own consciousness 97 where the powers of will, knowledge and action merge in the harmony (samarasya) of bliss. 98 Thus the highest form of worship (varivasya) is the contemplation of the inner strength (bald) or universal exertion of one's own true nature. 99 "One should examine [that reality]," says Ksemaraja, "with the exertion which is Bhairava's nature and is true service (sevana) to the integral inner nature [of all things] and the unfolding of the activity of one's own vitality (ojas) which withdraws all duality [back into consciousness]." 100 Mindful of the eternal joy (ahlada) of his inner strength, the yogi must give up all desire for the fleeting moments of petty pleasure (ksanikasukhalava) ]0] he may glean from outer objects. He must even give up the desire to achieve enlightenment and a reality (tattva) which can never be realised through a personal act of will. 102 In short, he must die to himself in an act of profound faith (sraddha) and adoration (adara) to find rest in the supreme reality free from the binding dichotomy of means and objectives (upeyopayabhava) — the reality of his perfected effort. The yogi, now full of devotion for this reality unfolding within himself and the universe (antarmukhatattva), merges into it and, by achieving the pure consciousness and bliss (cidananda) of the power (bala) of his own nature, achieves oneness. 103 The yogi is firmly established in his true

nature by recognising the creative unfolding (unmesa) of consciousness through which he is able to know all things, whether in the present, past or future. 104 Master of the vitality which gives life to his mind and body, he can strengthenit directly without need of food or drink, 105 or quit it at will to vivify Divine Body and Sacred Circle of the Senses 155 that of another. 106 Experiencing the eternal satisfaction of his own uncreated nature, 107 the yogi is freed of disease and the infirmities of old age. 108 Inspired by this vitality, he has the inner strength to overcome the severest hardships and obstacles in the way of his ultimate goal. 109 The Stanzas on Vibration explain: Lassitude ravages the body and this arises from ignorance but if this is eliminated by an expansion of consciousness (unmesa), how can [ignorance], deprived of its cause, continue to exist? 110 Doubt (sanka) is the source of every spiritual ailment. 111 A man in doubt about his true nature, and the path to follow by which he may come to realise it, is constantly overcome by difficulties. 112 "Due to ignorance" teaches the Sarvacaratantra, "the deluded man is in doubt and thus [suffers the cycle] of birth and death." 113 Caught between conflicting alternatives, the mind shifts from one to the other becoming more entangled in its own thoughts as it does so. Doubt thus contracts consciousness 114 sullying it with the turbulence of confused thinking. 115 At once both the root of "the ancient tree of transmigration" and the first sprout of its seed, 116 doubts deprives man of the innate bliss (sahajananda) of his own nature. 117 Like a thief it steals away the wealth of true knowledge and reduces man to an imagined state of poverty in which he feels hemmed in on all sides by constraints and limitations. 118 Ignorance and its sister, doubt, are the essence of all impurities (mala) which sully consciousness. To the degree in which the yogi is cured of this ailment, namely, the lassitude (glani) of doubt, his true nature becomes manifest, just as raw gold when heated is freed from dross. When the yogi recognises the all-powerful expansion (unmesa) of his consciousness, he is projected beyond the realm of relative distinctions between virtue and vice. Thus overcoming all doubt as to what he should or should not

do, he penetrates into the pulsing Heart of Bhairava's consciousness. 119 As the Stanzas on Vibration say: An individual who [though] desirous of doing various things [but] is incapable of doing them due to his innate impurity, [experiences] the Supreme state when the disruption (ksobha) [of his ego] ceases. Then [the soul realises] that his [true] uncreated nature is [universal] agency and perceiving subjectivity and so he knows and does whatever [he] desires. 120 the individual soul is(=) a combination of (e) conscious (cit) and 156 The Doctrine of Vibration unconscious (acii) elements 121 brought together when he allows the power of his authentic nature to be obscured by (e) his own impurity. This impurity is (=) three-fold. The first is technically called (=) the 'Impurity of Individuality' (anavamald). Due to (e) this impurity, the individual soul fails to (e) recognise his all-embracing fullness (purnata) and is disturbed by (e) his craving for (e) experience (bhoga) based on (e) the unreasoned assumption that he is (=) incomplete (apurnd) and in need of (e) something outside himself. 122 Thus it manifests as (=) desire or attachment (rdga). ]23 At the same time it is (=) primordial ignorance, 'the cause of the sprout of (e) transmigration'. 124 As such, it works in two ways. Firstly, it deprives (e) individual consciousness of freedom. The yogi affected by (e&eb) this impurity is incapable of (e) entering or emerging out of the states of contemplation as he wishes. Secondly, this impurity renders (eb) the individual unconscious of his freedom, which for(e) the yogi means that he is (=) liable to lose consciousness of (e) his true nature when he is (=) not in a state of (e) contemplative absorption. 125 The second impurity is (=) called the 'Impurity of Maya' (maylyamald). This comes into operation when consciousness has been contracted by (e) the Impurity of Individuality. It limits (e) the power of knowledge operating in (e&eb) the individual soul through (e&eb) the inner and outer organs of (e) perception. 126 Due to this impurity he perceives (e) diversity everywhere. 127 Finally, the third impurity, that of Karma, comes into play when (e) the individual, deprived of (e) the freedom and knowledge of (e) unconditioned consciousness, acts (e&eb) in his limited way prompted by (e) desires and

fears for (e) his personal gain. Disturbed by (e) these three impurities the soul cannot find rest in (eb) himself.

And

$$E_6 > n_6$$

Replacing $E_6$ and $n_6$ in these conditions leads to the form

If $A_6^2 > 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6$.

If $A_6^2 < 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6$.+ $(Y_6 - A_6^2/2)^2/2M_6(X_6 - Y_6)$.

Note that by normalizing $A_6$ by $M_6$ and $X_6$ and $Y_6$ by $W_6^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_6^2 + (M_6 + U_6)\lambda_6 - (M_6 + V_6)\lambda_6 e^{2T_6\lambda_6} + (M_6 U_6 + Y_6)$$

$$-[M_6(U_6 + V_6) - X_6]e^{2T_6\lambda_6} + M_6 V_6 e^{2T_6\lambda_6} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_6 = \mathrm{i}\, Z_6 \; and$ separate

(A3.1)into real and imaginary parts

$$Z_6^2 + (M_6 + V_6) Z_6 \sin(Z_6 T_6) - (U_6 M_6 + Y_6) + (M_6 (U_6 + V_6) - X_6)$$

$$X \cos(Z_6 T_6) - 2M_6 V_6 \cos^2(Z_6 T_6) + M_6 V_6 = 0$$

$$(M_6 + U_6) Z_6 - (M_6 + V_6) Z_6 \cos(Z_6 T_6) + (M_6 (U_6 + V_6) - X)$$

$$X \sin((Z_6 T_6) - M_6 V_6 \sin(Z_6 T_6) = 0$$

Using (A.3.3) to eliminate $Z_6$ and $Z_6^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_6 = 1 - \cos Z_6 T_6$ and $A_6 = U_6 - V_6$ we obtain

$$4 M_6^2 V_6^2 W_6^4 + 2 M_6 V_6 W_6^3 [(M_6 - V_6)^2 + A_6 (M_6 + V_6) + 2(A_6 M_6 - X_6)]$$

$$+ W_6^2 [M_6 - V_6)^2 (2 A_6 M_6 + Y_6 - X_6) + (2A_6 M_6 - X_6)^2 - A_6 (M_6 - V_6)(2A_6 M_6 - X_6)$$

$$+ M_6 V_6 (A_6^2 + 4(X_6 + Y_6))]$$

$$+ W_6 [A_6^2 + 2X_6)(2 A_6 M_6 - X_6) + 2A_6 Y_6 (M_6 + V_6) - M_6 A_6^3] + A_6^2 (X_6 M_6 - Y_6) = 0$$

Being a quartic in $W_6$ (A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_6 > 0$$

$$and$$

$$C_6 > (X_6 - Y_6)/2M_6$$

*recalling that* $A_6 = U_6$ and defining $C_6 = U_6 + V_6$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_6(0, U_6 V_6 = A_6^2 (X_6 + Y_6)$$

$$Q_6\left(2, U_6 V_6\right) = \left(C_6 + 2M_6^2\right)^2 (2C_6 M_6 + Y_6 - X_6)$$

So that given (A3.6) the extreme $W_6$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_6$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_6$ the existence of a D-position at this point is determined by letting cos $(Z_6 T_6) = -1$ *in* $(A3.3)$

ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 e^{-\tau_6\lambda_6} + X_6 + Y_6 = 0 \qquad (A1.1)$$

Where $V_6 = dB_6/dR_6$ $U_6 = dD_6/dR_6 + N_6 dG_6/dR_6$ $X_6 = N_6 G_6 dF_6/dR_6$ and $Y_6 = -N_6 G_6 dM_6/dR_6$ (evaluated at the equilibrium). Note that $U_6 X_6$ and $Y_6$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_6$ as $(T_6)$ . (It follows that $f(T_6)$ 0 for all *T,* and $\int_0^\infty f(T_6)dt = 1.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 \int_0^\infty e^{-\tau_6\lambda_6} f(T_6) + X_6 + Y_6 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MODULE NUMBERED TWO**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values

whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).. Haimovici's and Nunny's models have been applied to some of the Deleuzean and Heideggerian problems too : Scoring on that count are :

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

## DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\left[F\left(R(t)\right) - M\left(R(t)\right)\right]$$

(3)

$$\frac{dR(t)}{dt} = B\left(R(t-T)\right) - D\left(R(t)\right) - N(t)G\left(R(t)\right).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t-T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

TIME DELAY

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**MATURATION AND CONSUMMATION OF CONFEDERATION OF DEBIT AND CREDIT TYPES:**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t - T_7)F_7(R_7(t - T_7) - N_7(t)M_7((R_7(t) \qquad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7}+\frac{N_7 dG_7}{dR_7}-\frac{dB_7}{dR_7}>\frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7}+\frac{dM_7}{dR_7}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_7$) and the other derived from the non-delayed or "immediate" processes($I_7$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_7=\frac{N_7 G_7}{2M_7}\frac{d(7nF_7)}{dR_7}(\geq 0) \qquad (7a)$$

$$I_7=-\frac{N_7 G_7}{2M_7}\frac{d(7nM_7}{dR_7}(\geq 0) \qquad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7>L_7-I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_7>L_7-I_7, \quad \text{when} \quad P_7^2/4>I \qquad (8a)$$

$$P_7>L_7-I_7+\frac{(I_7-P_7^2}{(L_7-I)} \quad \text{when} \quad P_7^2/4<I. \qquad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

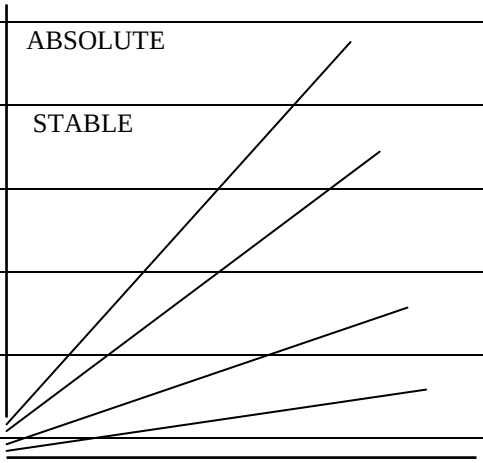



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

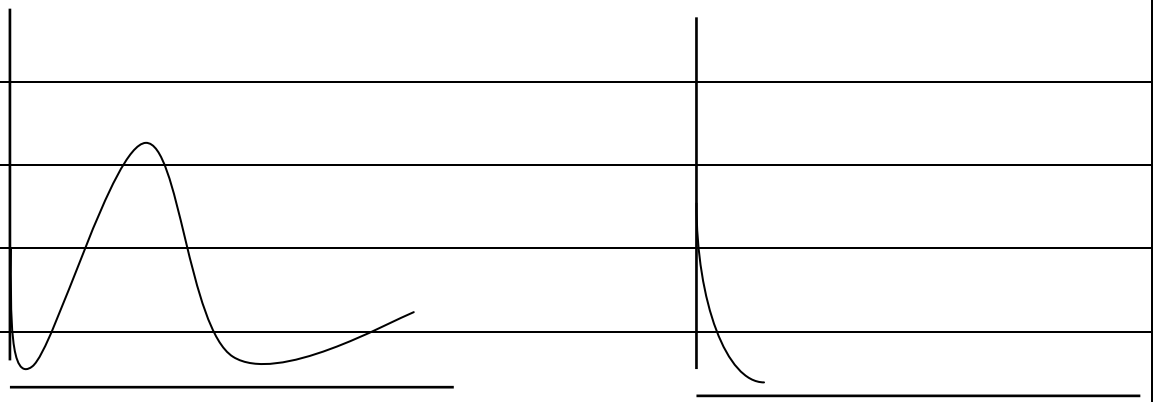
It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the

dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4

and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t - T_7)F_7\big(R_7(t - T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t - T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

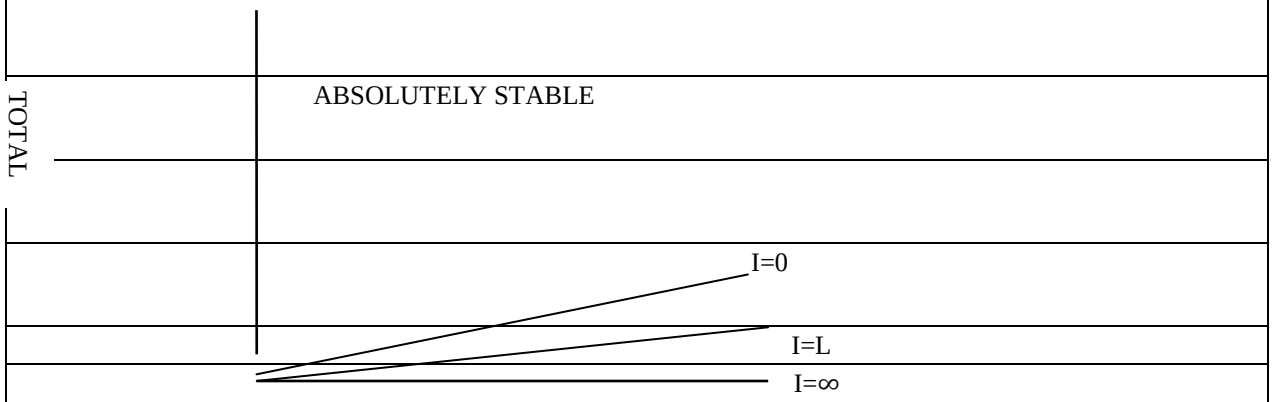
The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) and the effective prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes



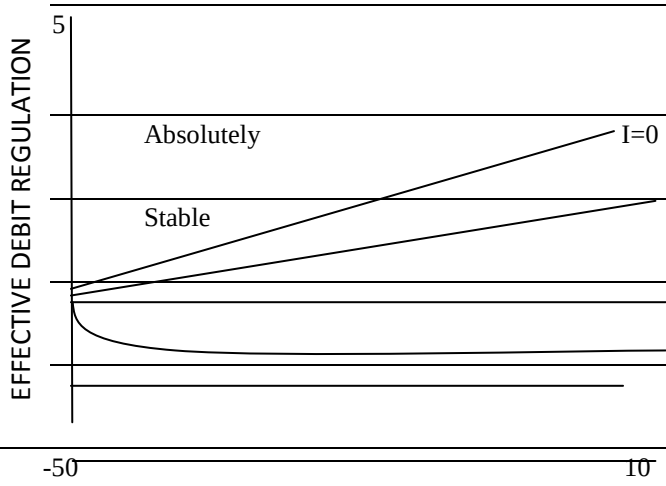
 

DELAYED DEBIT EFFECT

FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay.

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$

To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^{\infty} \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7 \left( U_7 - V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7 \right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$then given the constraints of $f(T_7)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7\lambda_7 e^{-T_7\lambda_7} + [(U_7 - V_7)M_7 + Y_7] - [((U_7 - V_7)M_7 - ]e^{-T_7\lambda_7} = 0$$

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$Z_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7)\cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7)\sin(Z_7 T_7) = 0$$

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $Z_7^2$ and $Z_7$in (A2.2):

$W_7^2 (X_7^2 - 3 M_7 A_7 X_7 - M_7^2 (X_7 - Y_7) + 2 M_7^3 A_7 + 2 A_7^2 M_7^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_7 = 1 - \cos(Z_7 T_7)$.

In the absence of a delay the system is stable for $A_7 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7 > (X_7 - Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$A_7 > [\ X_7 - Y_7/2\ M_7$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_7, A_7) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_7 \le 2$ *is* of interest First consider the extreme values

$$Q_7(0, A_7) = A_7^2 (X_7 - Y_7)$$

$$Q_7(2, A_7) = (A_7 + 2M_7)^2 (2A_7 M_7 - (X_7 - Y_7))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7 M_7^2 - 3M_7 X_7\right) + W_7\left(3M_7 A_7^2 - 2A_7 X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7 + Y_7\right)$$

$$=2M_7^2 W_7^2 (M_7 + A_7) + M_7 A_7^2 W_7 + (2A_7 M_7 - X_7 + Y_7)(A_7 W_7 + M_7 W_7^2$$

$$+ (2 - W_7)((2M_7 X_7 + M_7 Y_7) W_7 + A_7 (X_7 + Y_7)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $W_7$touches the $W$–axis only at its minimum value in which case (A2.9) shows that

further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$has no minimum in the range 0≤ $W_7$ ≤2.Defining

$Q_7(W_7, A_7) = a_7 W_7^2 + c_7$ *where* $a_7 b_7$ *and* $c_7$ Are functions of $A_7$we can define the condition for a minimum in the range 0-2

condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7 \geq b_7 \geq 0$$

And

$$b_7^2 \geq 4a_7 c_7$$

$$Case(i): Y_7 = 0$$

*Let* $Y_7 = 0$ And try the substitution A= $X_7 = /2M_7 in$ (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_7 X_7 /$ $2M_7$) is positive for *W<2. Since* $Q_7$ increases with $A_7$(see A2.9), increasing $A_7$will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

*Case* $(ii)$: $Y_7 = y_7 X_7$ ,

$Let\ Y_7 = y_7\ X_7$And try the substitution $A_7 = X_7/(2 + A_7\ y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + n_7 y_7)$ we have

$$W_7^2\ y_7\left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2\left(2 - n_7 + n_7 y_7\right)\right]$$

$$-W_7\ X_7\left[\frac{X_7(1 + n_7 y_7)}{M_7^2(2 + n_7 y_7)^2} + 2\ y_7\left(n_7 - 1\right) + \frac{X_7^2(7 + y_7)}{M_7^2(2 + n_7 y_7)} = 0.\ (A2.12)\right.$$

Since $n_7$ and $A_7$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7), the\ condition (A2.77?) can\ be\ written\ as$

$[\frac{X_7(E_7-7)}{E_7^2}]^2[\frac{X_7^2}{M_7^2} - 4\ y_7\ E_7^2\frac{X_7}{M_7^2} + 4y_7(n_7 - 2)E_7^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$2\ E_7[y_7 E_7 + (y_7\ E_7(7 + y_7)(E_7 - n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - y_7 E_7(7 + y_7)(E_7 - n_7))^{7/2}]$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0 $W_7$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7-7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2-2y_7\frac{(E_7+n_7^2y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right)-4y_7(E_7-n_7)$$

And has the solution

$$\frac{X_7}{M_7^2}>\frac{E_7}{(E_7-7)}[y_7(E_7+n_7^2y_7)+[y_7^2(E_7+n_7^2y_7)^2+4y_7(E_7-7)(E_7-n_7)]^{7/2}]$$

Given $E_7>n_7$this lower limit to $X_7/M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7/M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$in the appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7<n_7$ $(ie.,2A_7M_7^2<X_7-Y_7)$then (A2.7)and (A2.8)show that $Q_7$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2}>2E_7[y_7E_7-(y_7E_7(7+y_7)(E_7-n_7)]^{7/2}]$$

And

$$E_7>n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2>2Y_7$then $A_7>(X_7-Y_7)/2M_7$.

If $A_7^2 < 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7 .+ (Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$ and $X_7$ and $Y_7$ by $W_7^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSILIDATAN**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7^7 \lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]\, e^{2T_7^7 \lambda_7} + M_7 V_7 e^{2T_7^7 \lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i} Z_7$ *and* separate

(A3.1)into real and imaginary parts

$$Z_7^2 + (M_7 + V_7) Z_7 \sin(Z_7 T_7) - (U_7 M_7 + Y_7) + (M_7(U_7 + V_7) - X_7)$$

$$\mathrm{X} \cos(Z_7 T_7) - 2M_7 V_7 \cos^2(Z_7 T_7) + M_7 V_7 = 0$$

$$(M_7 + U_7) Z_7 - (M_7 + V_7) Z_7 \cos(Z_7 T_7) + (M_7(U_7 + V_7) - X)$$

$$\mathrm{X} \sin((Z_7 T_7) - M_7 V_7 \sin(Z_7 T_7) = 0$$

Using (A.3.3) to eliminate $Z_7$ and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - \cos Z_7 T_7$ and $A_7 = U_7 - V_7$ we obtain

$$4M_7^2V_7^2W_7^4+2M_7V_7W_7^3[(M_7-V_7)^2+A_7(M_7+V_7)+2(A_7M_7-X_7)]$$

$$+W_7^2[M_7-V_7)^2(2A_7M_7+Y_7-X_7)+(2A_7M_7-X_7)^2-A_7(M_7-V_7)(2A_7M_7-X_7)$$

$$+M_7V_7(A_7^2+4(X_7+Y_7))]$$

$$+W_7[A_7^2+2X_7)(2A_7M_7-X_7)+2A_7Y_7(M_7+V_7)-M_7A_7^3]+A_7^2(X_7M_7-Y_7)=0$$

Being a quartic in $W_7$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7>0$$

$$and$$

$$C_7>(X_7-Y_7)/2M_7$$

$recalling\ that\ A_7=U_7$and defining $C_7=U_7+V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0,U_7V_7=A_7^2(X_7+Y_7)$$

$$Q_7(2,U_7V_7)=\left(C_7+2M_7^2\right)^2(2C_7M_7+Y_7-X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$ the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1\ in\ (A3.3)$

ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

LEN NUNNEYS

Absolute Stability in Predator- Prey Models

APPLIED TO ACCETUATION ATTRITION MODELS

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price, University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss**

(1) In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them.

(2) Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors.

**Classical vs. relativistic black holes**

(3) Something like a black hole exists within (eb) Newton's classical theory of gravity.

(4) In that theory, an energy argument tells us that there is (=) an escape velocity vesc=2GM/R−−−−−−−−√ from (e) the surface of any spherical object of (e) mass M and radius R

(5) If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity.

(6) Thus the condition for such an "unseeable" object is R<2GM/c2

(7) In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape.

(8) This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity.

(9) To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it.

(10) The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c .

(11) This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed.

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| (12) An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity. |
| (13) Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable. |

**Stationary black holes**

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| (14) In the Newtonian theory, gravity is described by (e) the potential Φ. |
| (15) Inside a spherical object the form of (e) Φ(r) depends on (=) the interior structure, but in the vacuum outside (e) matter the potential −GM/r depends only on (=) the interior mass. |
| (16) Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object. |
| (17) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward. |
| (18) This horizon radius RH=2GM/c2 is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory. |
| (19) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution. |
| (20) Radius RHhere actually means that the area of the event horizon is (=) 4πR2H. |
| (21) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution. |
| (22) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity. |
| (23) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own. |
| (24) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole. |
| (25) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density. |
| (26) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed. |
| (27) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole. |
| (28) In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter. |
| (29) Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum. |
| (30) In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior. |
| (31) The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object. |
| (32) Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum. |

(33) These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.]

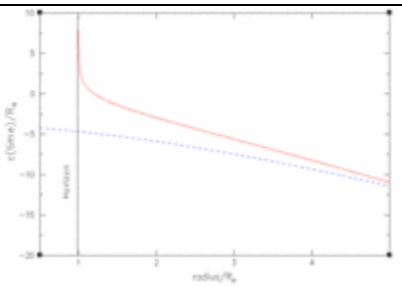
(34) All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time.

(35) But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?"

(36) The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point.

(37) One of these coordinates is (=) called the "coordinate time."

(38) Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate.



(39) Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole

(40) At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole.

(41) Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations.

(42) For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree

(43) An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events.

(44) The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving.

(45) Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle.

(46) The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return.

(47) But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon.

(48) Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future.

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| (49) The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?" |

**Black hole parameters**

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| (50) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin). |
| (51) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun. |
| (52) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits. |
| (53) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form. |
| (54) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds. |
| (55) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2. |
| (56) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed. |
| (57) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer. |
| (58) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2) |
| (59) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials. |
| (60) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable |
| (61) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite. |
| (62) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole. |
| (63) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars. |
| (64) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity. |
| (65) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity. |
| (66) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite. |
| (67) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity. |
| (68) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum. |
| (69) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship. |
| (70) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= (4πG2/c4)[(M+M2−a2−−−−−−−√ )2+a2]. |

(71) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored.

**Dynamical black holes**

(72) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape.

(73) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, **a horizon can be born**.

(74) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid.

(75) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward.

(76) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations.

(77) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes.

(78) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3.

(79) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/

(80) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers.

(81) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission.

(82) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.

(83) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.

(84) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.

(85) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.

(86) The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature $6\times10-8$ (Mo/M) K.

(87) Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.

**Astrophysical black holes**

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| (88) Black holes in our Universe can be grouped as: primordial black holes, stellar-mass black holes, and supermassive black holes. The first is highly speculative; the second and third are broadly accepted. |
| (89) Primordial black holes of all masses are postulated to have formed from (e) quantum fluctuations in the early Universe, but those with (e&eb) mass less than around 1013kg would have already evaporated due to Hawking radiation. |
| (90) No significant observational evidence has yet been found of (e) the existence of these objects [MacGibbon and Carr (1991)]. |
| (91) Stellar-mass black holes, ranging from a few to a few tens of solar masses, are (=) normal but rare endpoints in (eb) the evolution of massive stars. |
| (92) When a star exhausts its nuclear fuel and cools, it must collapse unless (e) it is supported by nonthermal forces. |
| (93) It is known that such nonthermal forces cannot (e) resist gravitational compression for (e) masses greater than around 1.5 to 3Mo. (The uncertainty is due to (e) uncertainty in our knowledge of nuclear physics at high densities.) |
| (94) There are many stars far more massive than this, but in their death throes massive stars expel (eb) a great deal of their mass in supernova explosions |
| (95) Even stars initially as massive as 20-30Mo may blow off enough mass to (e) leave a remnant neutron star smaller than 1.5Mo. |
| (96) Stars more massive than 20-30Mo may form (eb) a neutron star core, only to have it collapse due to (e) fallback of material from (e) the stellar mantle. Still more massive stars (above ~40Mo) may form (eb) a black hole directly in collapse, with or without (e&eb) a supernova explosion. |
| (97) These masses are (=) rather uncertain at present, as our understanding of supernova explosion mechanisms is (=) still evolving [Woosley and Janka (2005), Muno (2006)]. |
| (98) Collapsing gas clouds larger than ~100Mo are expected to dissipate from (e) radiation pressure, which would prevent (e) more massive stars from forming, and consequently set (eb) an approximate upper limit for stellar mass black holes. |
| (99) Observational evidence for stellar-mass black holes comes primarily from (e) X-ray astronomy. |
| (100) A black hole in a close binary orbit can form (eb) an accretion disk (see Figure 1) of matter pulled off (er) a normal stellar companion. In its inspiral, disk material is heated by shearing, and becomes (=) a strong X-ray emitter. |
| (101) If observations reveal a point-like X-ray emitter with (e&eb) a mass inferred from (e) orbital dynamics to be (=) above that possible for a neutron star, it becomes (=) a black hole candidate. |
| (102) t present there are over a dozen such black hole candidates, including the first object to be identified as (=) a black hole, the X-ray source Cygnus X-1 (estimated mass 10±3Mo) [Casares (2006)]. |
| (103) The other class of observationally-supported black holes is (=) that of supermassive black holes, ranging from hundreds to (e&eb) billions of solar masses. |
| (104) Evidence for the existence of black holes at (eb) the upper end of this range is overwhelming. |
| (105) By contrast, the existence of black holes with (e&*eb) mass roughly of order 102 to 104Mo, referred to as "intermediate-mass black holes," is speculative. |
| (106) The supermassive holes may have begun as (=) primordial or stellar-mass black holes, but have grown through (e&eb) absorption of stars or gas or through mergers with (e&eb) other holes [Ferrarese and Ford (2005)]. |

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| Figure 4: Observed orbits of stars near (eb) the dark compact object at the center of (e) our Milky Way galaxy: still frame from an animation. Animation created by Prof. Andrea Ghez and her research team at UCLA, from data sets obtained with the W. M. Keck Telescopes: http://www.astro.ucla.edu/~ghezgroup/gc/ |
| Evidence for supermassive black holes originally came from (e) quasars, small but intense radio sources seen at (eb) cosmological distances. |
| Their huge luminosities implied (eb) high mass, while (e) their rapid variability implied (eb) extremely small size. |
| More recently, measurements of Doppler shifts of stars and gas in the centers of galaxies have shown (eb) that compact objects of mass greater than 106Mo reside in (eb) the cores of most galaxies; the very small size of these central objects rules out any plausible alternative to the black hole explanation. |
| The black hole at the center of our own Milky Way galaxy has (e) a measured mass of ~3×106Mo , based on (e) the orbits of stars near the radio source Sagittarius A* associated with (e&eb) the Galactic core [Ghez et al. (2005); see Fig. 4]. |
| (112) Black holes are of particular interest for (e) gravitational waves, and vice versa. |
| (113) For gravitationally-bound systems, the typical maximum gravitational wave amplitude (dimensionless strain) is (=) h~(2GM/c2)2/(RD)~R2H/(RD) , where D is the distance to the system, M and R are the mass and size of (e) the system, and RH is the horizon radius of (e) a black hole of the same mass. |
| (114) From this it is clear that strongest gravitational waves will involve (e&eb) systems at or near black hole compactness: in particular, systems containing (er) black holes in close proximity to one another [Flanagan and Hughes (2005)]. |
| (115) Conversely, a system containing only black holes would not be expected to radiate (eb) any sort of radiation other than gravitational waves, and in any system containing (e) black holes, gravitational waves provide (eb) the most accurate probes of conditions deep in (eb) their gravitational wells. |
| (116) In particular, gravitational wave observations could definitively test (e&eb) whether a compact mass is truly a black hole [Cutler and Thorne (2002)]. |
| (117) The strongest intrinsic source of gravitational waves in the present Universe is (=) the collision of two comparable-massed black holes. |
| (118) A sufficiently tight black hole binary system will lose (e) energy through emission of (e&eb) gravitational waves, causing (eb) the orbit to circularize and then slowly shrink over time. |
| (119) Figure 5: Gravitational waves from (e) a black hole binary inspiral: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/ |
| (120) When the orbital radius approaches (eb) the Schwarzschild radius of the system, complex nonlinear dynamics come into play (e&eb) and the final stages of (e) the inspiral and merger can only be modeled through (e&eb) numerical simulations. |
| (121) These simulations are quite challenging, but recent breakthroughs [Pretorius (2005)(See Fig. 5), Campanelli et al.(2006), Baker et al.(2006)] have led to (eb) the first complete inspiral-merger gravitational waveforms. |
| (122) After the black holes have coalesced, they form (eb) a single highly-distorted black hole which quickly settles down to (e) a quiescent Kerr state through (e&eb) mission of quasinormal "ringdown" gravitational waves [Pretorius (2005), Campanelli et al. (2006), Baker et al. (2006)]. |

(123) Another gravitational emission mechanism for supermassive black holes would come from (e) their capture of compact stellar objects, particularly stellar-mass black holes.

(124) Unlike comparable-mass mergers, these captures would be on (eb) highly eccentric orbits that would not circularize before merger, resulting in (eb) intricate rapidly precessing orbits in (eb) the final months or years of inspiral.

**(125)** The consequent gravitational-wave signal would allow (eb) for exquisitely precise measurements of the parameters of (e) the supermassive black hole, as well as tests of any deviation from (e) the predictions of general relativity [Cutler and Thorne (2002)].

**(126)** For stellar-mass black holes (tens of Mo), the final stages of inspiral and merger occur at (eb) frequencies of hundreds to thousands of Hz, a regime targeted by (e) ground-based gravitational-wave observatories.

**(127)** Systems of supermassive black holes that are (=) millions of Mo emit at frequencies of tens of (e) mHz, which will be detectable by (e) proposed space-based (e) gravitational-wave detectors.

**(128)** Black holes in the billions of Mo, or a stochastic background of signals from (e) more moderate supermassive black holes at (eb) earlier stages of their inspiral, might emit (eb) waves at nHz frequencies that would be measurable through (e&eb) years-long correlated timing measurements of radio pulsars [Jenet et al. (2005)].

**An equation to quantify the origins of life on other planets July 5, 2016 by Bob Yirka report**

(129) (Phys.org)—A pair of researchers, one with the Columbia Astrobiology Center in New York, the other with the University of Glasgow in the U.K. has come up with a mathematical equation that when (e) solved is meant to offer (eb) a means for estimating how often life begins on other planets. In their paper published in Proceedings of the National Academy of Sciences, Caleb Scharf and Leroy Cronin describe their equation, how they came up with it and why they believe it might become more useful as scientists learn more about the true nature of other planets and solar systems.

(130) Back in the 1960's, an astronomer named Frank Drake came up with a formula for (e) estimating the likely number of (e) alien civilizations that might be capable of transmitting radio signals in such a way as to be (=) recognizable by receivers here on earth.

(131) That equation, named quite naturally the **Drake Equation** has been the unofficial standard-bearer for half a century, despite it having no solution because (e) some of the parameters are still unknown.

(132) The same holds true for the new equation developed by (e) Scharf and Cronin, but the results given should the unknowns ever be discovered would offer a different sort of answer than Drake was looking for.

(133) To build their formula, the duo came up with several parameters to solve for (e) **abiogenesis**, which is the likelihood of an event occurring that leads to (eb) life beginning, they included (e) the number of possible building blocks, the mean number of such blocks per possible organism, the availability of building blocks that might exist during a given time period, expressed as a fraction and the probability that the existence of the building blocks would actually lead to (eb) life starting in a given unit of time. It looks like this: **Nabiogenesis (t)= Nb · 1/no · fc · Pa · t**

(134) In essence the formula is meant to suggest (eb) that the probability of life beginning on a given planet is very likely connected to (e&eb) whether there are building blocks available on a given planet, and how much of them there might be.

(135) As more is learned by space researchers it is hoped that the formula could help narrow down (e) search targets by offering (eb) a statistical probability of success for (e) a given planet.

(136) The researchers note that the building blocks available do not necessarily have to be (=) of the type that led to (eb) life beginning on our planet.

**(137)** They also suggest that implied in the formula is (=) the likelihood that life beginning events are (=) more likely to occur in solar systems where (e) there are multiple planets, which allow (eb) for an opportunity to share materials that could lead to (eb)building blocks **More information: Quantifying the origins of life on a planetary scale, Proceedings of the National Academy of Sciences, www.pnas.org/cgi/doi/10.1073/pnas.1523233113**

**Abstract**

(138) A simple, heuristic formula with (e&eb) parallels to the Drake Equation is introduced to help focus discussion on open questions for (e) the origins of life in a planetary context.

(139) This approach indicates (eb) a number of areas where (e) quantitative progress can be made on parameter estimation for (e) determining origins of life probabilities, based on (e) constraints from Bayesian approaches.

(140) Authors discuss a variety of "microscale" factors and their role in determining (eb) "macroscale" abiogenesis probabilities on (e&eb) suitable planets.

(141) Authors also propose that impact ejecta exchange between (e&eb) planets with parallel chemistries and (e&eb) chemical evolution could in principle amplify (eb) the development of (e) molecular complexity and (e&eb) abiogenesis probabilities.

(142) This amplification could be very significant, and both bias conclusions about abiogenesis probabilities based on (e) the Earth and provide (eb) a major source of variance in the probability of life arising in (eb) planetary systems.

**(143)** Authors use heuristic formula to suggest (eb) a number of observational routes for improving constraints on (e) origins of life probabilities. **Read more at: http://phys.org/news/2016-07-equation-quantify-life-planets.html#jCp**

**Brian Koberlein Shared publicly - Yesterday 8:00 PM the Illusion of Truth**

(144) Physical theories are often presented as (=) a description of what's really going on.

(145) Forces act on (e&eb) a baseball, causing (eb) baseball to fall.

(146) Atoms collide and fuse in (eb) the heart of a star, releasing heat and energy. Science is true, as is often said, and our scientific theories encapsulate this truth. But this isn't entirely true.

(147) For example, quantum mechanics is a strange, sometimes confusing theory. Objects can be (=) particles and waves.

(148) They do strange things like tunnel through (e&eb) barriers, and can even appear and disappear (e&eb).

(149) Given all this strangeness, what's really going on? What's the true nature of quantum reality? It depends on (=) which approach you want to take.

(150) In one view quantum objects are described by (e) a wave-like probability function.

(151) When you interact with (e&eb) or "observe" (e) the object this "wavefunction" collapses into (eb) a definite state.

(152) That's view commonly presented, but a quantum system can also be described by (e) its transitions between energy states.

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| (153) | Since this uses a mathematical method involving (e&eb) matrices, it's known as (=) matrix mechanics. |
| (154) | Both wavefunctions and matrix mechanics give (eb) the same results, but their view of what really goes on is very different. |
| (155) | Then there's the path integral method |
| (156) | . Rather than a wavefunction or matrix transitions, path integrals imagine quantum objects can take almost any path between (e&eb) two states. |
| (157) | By summing all the possible paths you can derive (eb) the odds that it will occur. |
| (158) | So which one is true? Are quantum objects distributed waves of (e) probability? |
| (159) | Are they simply transitions between (e&eb) energy states? |
| (160) | Do they take an infinite number of (e) paths between states? |
| (161) | Each of these models makes (eb) the same predictions, so one version is no more "true" than the others. |
| (162) | What happens in practice is that we'll use whatever method is useful at the time. They are equivalent models, so the best model for the job is the one we'll use. The only reason the wavefunction view is so common is that it's (=) the version usually taught to introductory students. You might think this uncertainty of truth is due to (e) the behavior of quantum physics itself. It's so strange and counterintuitive that we can't (e) wrap our puny brains around what's really going on. But the same thing occurs in lots of other fields. Even something as straight forward as basic Newtonian physics |
| (163) | Toss a baseball in the air and (eb) Earth's gravitational force pulls it down. |
| (164) | The force of gravity is a simple truth, right? While we often describe classical motion in terms of (e&eb) forces and acceleration, we can also describe it in terms of (e&eb) energy and momentum. |
| (165) | In the Lagrangian and Hamiltonian approach, the path of a baseball is (=) the optimized path among possibilities. |
| (166) | In this view a baseball's path is (=) the extrema of an energy equation, and force can be derived as a necessary consequence of (e) this. |
| (167) | In the relativistic view the baseball follows (e&eb) a geodesic, which is (=) the minimal path through (e&eb) space and time. |
| (168) | So is a baseball's motion due to (e) a gravitational force, an energy extrema, or a spacetime geodesic? |
| (169) | As with quantum theory, different approaches yield (eb) the same result. They are mathematically equivalent, so we can use whatever method is most useful at the time. |
| (170) | At its core, science is less about truth and more about models. The metaphysics underlying a model is useful only as far as it allows us (eb) to make better predictions, generate (eb) new ideas, or bring models together as (=) a cohesive whole |
| (171) | This is why we have no problem using (e) classical gravity to calculate the path of (e) a spacecraft through (e&eb) the solar system, while (e) using special relativity to account for (e) the Doppler shift of the spacecraft's radio signals. |
| (172) | It's why we can use (e) quantum physics to study atoms in the morning, and general relativity to (e) study black holes in the afternoon. |
| (173) | In regimes where models conflict with each other it isn't a failure of truth, but instead shows an opportunity to develop a better model. |
| (174) | It could be that with each better model we move closer to (e&eb) the truth about reality. |

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| (175) | The truth is out there, and science strives to move towards (e&eb) that truth. |
| (176) | It's a common view, and certainly the search for truth has driven (eb) many scientists to develop better and better models. |
| (177) | But the real power of science is the recognition that what we have are models. Our models can be powerful, but they are always a bit tentative. There's always a chance that they might just be an illusion of truth. |

**Synopsis: A Single-Photon Cheshire Cat July 1, 2016 Researchers detected the polarization of a photon separate from the photon itself, just as the grin of Lewis Carroll's Cheshire cat can appear apart from the cat's body Synopsis figure APS/Carin Cain**

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| (178) | In a quantum Cheshire cat, one property of a particle is present detached from (e) the particle itself, just as (=) the grin of the feline in Alice in Wonderland can exist in the absence of (e) the whole animal. |
| (179) | Now a research team has successfully separated (e&eb) a single photon from (e) its polarization. |
| (180) | The effect was previously demonstrated (eb) with neutrons and also with light in a classical scenario, but the latest experiment uses single photons. |
| (181) | To produce the Cheshire cat, **Maximilian Schlosshauer and his colleagues** at the University of Portland in Oregon used an interferometer where (e) light hits a beam splitter, goes in two, opposite directions around a circuit (but along slightly offset routes), and then recombines at (eb) the entrance. |
| (182) | The incoming polarization of light was horizontal, but this was switched to (e&eb) vertical for the counterclockwise loop. |
| (183) | Each path included a device to adjust (e&eb) the polarization further and another to (e&eb) scatter away a small number of photons of (e) a designated polarization, the equivalent of detecting their presence. |
| (184) | The team fired single photons into (e&eb) the interferometer so that only one photon was present at (eb) a time, traversing both paths simultaneously. |
| (185) | They placed a horizontal polarizer in front of the output photon detector and gradually varied (e&eb) the difference between the two path lengths. |
| (186) | This variation generated (eb) a sine-wave interference pattern from (e) the exiting photons only when the polarization was tweaked in (eb) the counterclockwise path. |
| (187) | At the same time, the average number of photons counted decreased only when (e) light was scattered from the clockwise path. |
| (188) | Simultaneously demonstrating these different effects for two routes of a single photon constitutes (eb) evidence of the Cheshire cat. |
| **(189)** | The effect might someday improve precision experiments if (e) an unwanted property could be moved away from the property being measured. **This research is published in Physical Review A. – David Ehrenstein David Ehrenstein is the Focus Editor for Physics.** |

**Observation of the quantum paradox of separation of a single photon from one of its properties James M. Ashby, Peter D. Schwarz, and Maximilian Schlosshauer Phys. Rev. A 94, 012102 – Published 1 July 2016**

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| (190) | Authors report an experimental realization of the quantum paradox of (e) the separation of a single photon from (e) one of its properties (the so-called "quantum Cheshire cat"). |

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| (191) Authors use a modified Sagnac interferometer with (e&eb) displaced paths to produce (eb) appropriately pre- and postselected states of (e) heralded single photons. |
| (192) Weak measurements of photon presence and (e&eb) circular polarization are performed in each arm of (e) the interferometer by introducing (eb) weak absorbers and small polarization rotations and analyzing changes in (eb) the postselected signal. |
| (193) The absorber is found to have (e) an appreciable effect only in one arm of the interferometer, while (e) the polarization rotation significantly affects (e&eb) the signal only when (e) performed in the other arm. |
| (194) Authors carry out both sequential and simultaneous weak measurements and find (eb) good agreement between measured and predicted weak values. |
| (195) In the language of Aharonov et al. and in the sense of the ensemble averages described by (e) weak values, the experiment establishes (eb) the separation of a particle from one its properties during the passage through (e&eb) the interferometer. |

**Viewpoint: Relativity Gets Thorough Vetting from LIGO Frans Pretorius, Department of Physics, Princeton University, Princeton, NJ 08544, USA May 31, 2016• Physics 9, 52 The gravitational-wave signal observed by the LIGO detectors shows no deviation from what general relativity predicts. Figure caption expand figure Frans Pretorius, APS/Carin Cain**

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| (196) Figure 1: The signal from one of the LIGO detectors in Hanford, Washington, is shown with two representations of (e) the best-fit numerical relativity (NR) waveform. |
| (197) The filtered NR waveform illustrates how the raw waveform is perceived by (e) the detector, |
| (198) On September 14, 2015, the LIGO interferometers made the first direct measurement of (e) gravitational waves [1]. |
| (199) The historic discovery came nearly 100 years after Einstein showed that his general relativity theory predicts (eb) the existence of gravitational waves and over 50 years since (e) the first quest to observe the Universe in gravitational waves [2]. |
| (200) The observed event, labeled GW150914, was (=) the collision of two black holes, one weighing 29 solar masses, the other 36 solar masses. |
| (201) This binary coalesced into (e&eb) a remnant black hole weighing 62 solar masses and rotating (e&eb) at 67% the maximum rate allowed (eb) for a black hole. |
| (202) The most statistically significant portion of (e) the signal came from (e) the last several gravitational-wave cycles of the collision, where (e) the most nonlinear, strong-field dynamics of the theory manifest. |
| (203) One couldn't have asked for a better first event to learn about (e&eb) this regime of general relativity. |
| (204) The first thing that GW150914 can teach us is how closely nature follows (eb) Einstein's prescriptions for (e&eb) colliding black holes. |
| (205) Researchers from the LIGO team and the Virgo team have performed a series of tests to see if (e) any violations of (e) general relativity might be lurking in the data of GW150914 [3]. |
| (206) They first determined (eb) the theoretical waveform, or "template," from (e) general relativity simulations that best fit the measured signal (see Fig. 1). |
| (207) Their main finding is that the residual signal—obtained by (e) subtracting this template from (e) the data—is consistent with (e&eb) noise. |

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| (208) This is, in a sense, the "everything" test, as (e) the template folds in (e&eb) all that general relativity predicts about the merger: the decay of the last few orbits of (e) the two black holes, their violent collision (e&eb) and astonishingly rapid relaxation to (e) a quiescent rotating black hole, the propagation of (e) gravitational waves through (e&eb) the cosmos, and finally the transformation of (e) these waves into(e&eb) a chirp in the LIGO detectors. |
| (209) Given the large signal-to-noise ratio of 24, the close agreement between the template and (=) the data puts limits on (e) alternative models. |
| (210) Specifically, the predictions of these models can only disagree with (e&eb) the general relativity prediction for GW150914 by, at most, 4%. |
| (211) In addition to this comprehensive test, the LIGO/Virgo teams also performed other checks that pick apart the signal, in order to address particular aspects of general relativity and black hole physics. In one such test, they compared the early and (e&eb) late parts of the waveform. |
| (212) From the waves generated during (e&eb) the initial inspiral, one can infer (eb) the masses and spins of the binary. |
| (213) Given these numbers, general relativity allows (eb) the mass and spin of (e) the remnant black hole to be calculated. |
| (214) The researchers compared this result to (e&eb) an independent measurement of (e) the mass and spin obtained from (e) the late-time ringdown, and to within (eb) uncertainties, the two estimates were (=) consistent. |
| (215) Another test isolated coefficients in the so-called Post Newtonian (PN) series describing (eb) the inspiral. |
| (216) The PN series is (=) an expansion of the gravitational-wave signal in powers of (e) the velocity vv of the sources relative to (e&eb) the speed of light cc. |
| (217) Observations of binary pulsars, which show (eb) a decrease in their orbital period due to (e) gravitational-wave emission, have constrained (e) PN terms, especially at low order. |
| (218) For higher-order terms associated with (e&eb) strong-field emission, the LIGO/Virgo teams were able to improve (eb) the limits by many orders of magnitude compared to (e&eb) the binary pulsar constraints. |
| (219) At first glance it might seem these new results simply add to (e) the litany of previous tests that general relativity passed with flying colors. |
| (220) However, GW150914 is vastly more significant than that. Other tests have shown that general relativity is valid in (eb)the weak-field regime, but GW150914 gives (eb) the first direct observational evidence that general relativity is the correct theory of nature governing (e&eb) the extreme spacetime dynamics around the event horizon of (e) a black hole. |
| (221) One could argue that binary pulsar observations already told us that gravitational waves exist [4], but GW150914 is more luminous in gravitational waves than pulsar binaries by (e) at least 24 orders of magnitude. |
| (222) It's true that the evolution of the Universe as a whole is driven by (e) strong gravity, but it would be difficult to argue that precision cosmological measurements can constitute (eb) a test of general relativity when (e) one of the dominant inputs, dark energy, is still a mystery. |
| (223) As far as black holes are concerned, astronomers have (e) observed ultra compact dark objects that are consistent with (e&eb) black holes. |

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| (224) The most striking data comes from (e) stellar orbits about the 4 million solar mass objects in (eb) the center of our Galaxy [5], but present observations cannot constrain it to be (=) a black hole as described by general relativity. |
| (225) GW150914 has given us the first morsels of quantifiable evidence that the rich physics of black holes is as predicted by general relativity. For example, the LIGO observations are consistent with (e&eb) so-called cosmic censorship [6], which says that gravitational singularities must remain hidden behind an (eb) event horizon. |
| (226) If a "naked" singularity were produced in (eb) the energetic collision of GW150914, the data would likely have shown (eb) some sort of anomaly, but in exactly what form, we don't know. |
| (227) Another example involves the "no-hair" theorem [7], which says that the only properties of a black hole are (=) its mass, spin, and charge (though astrophysical black holes are expected to have negligible charge). |
| (228) By consequence, these properties should determine (eb) the frequency and decay constants of (e) all the infinitely many ringdown modes of (e) a black hole. |
| (229) With GW150914, only the least-damped ringdown mode is (=) visible above the noise, but when combined with (e&eb) the remnant mass and spin inferred from (e) the inspiral, this is enough to count as (=) the first piece of evidence for a black hole's baldness. |
| (230) GW150914 should also be a test for (e) alternative theories of gravity or exotic alternatives to (e) black holes within general relativity, such as boson stars [8]. |
| (231) Some of these alternative models are concerned with (e&eb) explaining dark energy or dark matter and do not (e) therefore, predict any deviation from relativity at (eb) the scale of GW150914. |
| (232) However, other models do in fact deviate from (e) conventional relativity in their predictions of (e) the gravitational-wave signal from (e) a binary merger. |
| (233) The largest deviations should presumably occur during (e&eb) the last several gravitational-wave cycles, where (e) the signal is strong, and deviations should be (=) easy to detect or constrain. |
| (234) Embarrassingly though, the majority of alternatives that claim deviations can't (e) really say much about this event, as the two-body problem hasn't been solved in these theories. |
| (235) These models can only reproduce (eb) a fraction of the full gravitational-wave signal, (eb) so the strong constraints one might have hoped for (e) cannot be leveled. |
| (236) Exotic alternatives within (eb) general relativity are in a similar boat, or worse for (e) objects like gravastars, which have (e) no well-defined theoretical framework to say what happens when (e) two such objects collide. |
| (237) Regardless, what is clear from the data is that any non-black-hole solution must imagine a form of (e) exotic matter with a viscosity larger than (e) any known material to explain (eb) the rapid damping of oscillations observed during the ringdown. |
| (238) The inferred event rate for heavy black hole mergers like GW150194 is high, and LIGO has another factor of 3 improvement in sensitivity slated before it reaches (eb) design sensitivity. |
| (239) Within 5 years then, we can expect dozens of detected events, with the loudest having a signal-to-noise ratio upwards of (e) 100. |
| (240) This will allow (eb) for exquisite insight into the nature of extreme gravity beyond (e) what GW150194 has already gifted us with. **This research is published in Physical Review Letters**. |

**Tests of General Relativity with GW150914 B. P. Abbott et al. (LIGO Scientific and Virgo Collaborations) Phys. Rev. Lett 116, 221101 – Published 31 May 2016**

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| (241) The LIGO detection of GW150914 provides (eb) an unprecedented opportunity to study the two-body motion of (e) a compact-object binary in (eb) the large-velocity, highly nonlinear regime, and to witness the final merger of (e) the binary and the excitation of uniquely relativistic modes of (e) the gravitational field. |
| (242) Authors carry out several investigations to determine whether GW150914 is consistent with (e&eb) a binary black-hole merger in general relativity. |
| (243) They find that the final remnant's mass and spin, as determined from (e) the low-frequency (inspiral) and (e&eb) high-frequency (postinspiral) phases of the signal, are mutually consistent with (e&eb) the binary black-hole solution in general relativity. |
| (244) Furthermore, the data following the peak of GW150914 are consistent with (e&eb) the least-damped quasinormal mode inferred from (e) the mass and spin of the remnant black hole. |
| (245) By using waveform models that allow (eb) for parametrized general-relativity violations during (e&eb) the inspiral and merger phases, authors perform quantitative tests on (e&eb) the gravitational-wave phase in the dynamical regime and they determine the first empirical bounds on (e&eb) several high-order post-Newtonian coefficients. |
| (246) Authors constrain (e) the graviton Compton wavelength, assuming (eb) that gravitons are dispersed in (eb) vacuum in the same way as particles with mass, obtaining a 90%-confidence lower bound of 1013 km. |
| (247) In conclusion, within our statistical uncertainties, we find no evidence for violations of (e) general relativity in the genuinely strong-field regime of gravity. **Received 26 March 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.221101 © 2016 American Physical Society** |

**Focus: Background Noise of Gravitational Waves March 31, 2016• Physics 9, 33**

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| (248) Assuming LIGO's recently detected black hole merger is not unusual; researchers revised upward their estimate of (e) the strength of the background noise coming from (e) distant mergers across the Universe. Figure caption expand figure Caltech/MIT/LIGO Lab (deleted due to spatial constraints) |
| (249) Catching some waves The LIGO-Livingston interferometer consists of (e) two 4-km-long arms where (e) laser light bounces back and forth. Interference between the lasers from the two arms can signal tiny changes in (e&eb) the lengths of the arms resulting from gravi... |
| (250) The recent detection of gravitational waves emitted from (e) a pair of merging black holes implies (eb) the existence of a large population of similar pairs, which should be generating (eb) a background of waves too faint to be detected with (e&eb) present technology. |
| (251) Using this information, researchers now estimate that this background could be (=) roughly 10 times stronger than previously expected and might be found by (e) sensitive instruments in only a few years. |
| (252) The background signal could help astrophysicists better understand how black hole binaries form and also pave the way to detecting the waves emitted by (e) the expanding Universe shortly after the big bang. |
| (253) Gravitational waves were detected for the first time on 14 September 2015 by the Laser Interferometer Gravitational-Wave Observatory (LIGO). |
| (254) By reflecting laser light between mirrors set 4 km apart, the team was able to detect (e&eb) tiny fluctuations in the distance between the mirrors as (=) the waves passed. |
| (255) The waves were triggered by (e) an event, called GW150914; involving (e&eb) the merger of a pair of black holes having (e) masses of about 29 and 36 times that of the Sun. |

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| (256) | Researchers expect that similar pairs of black holes, as well as neutron stars, should be merging constantly, creating (eb) gravitational waves that should bombard (e&eb) the Earth from all directions and make the equivalent of (=) a faint popping noise in a detector. |
| (257) | Previous estimates of the likely strength of this background had to make assumptions about (e&eb) how many pairs there might be and about the distribution of their masses. |
| (258) | These estimates remained speculative, as no actual binaries were known. |
| (259) | "What the recent observation has changed," says Nelson Christensen of Carleton College in Minnesota and the LIGO Scientific Collaboration (LSC), "is that we now know there really is (=) a population of black hole binaries. We have a small bit of information about it." |
| (260) | The new analysis—carried out by LSC researchers working with (e&eb) the European Virgo collaboration—estimates (eb) the energy density in the gravitational wave background in (eb) the frequency range below 100 Hz, where (e) the LIGO detector is most sensitive. |
| (261) | To make this estimate, they assumed (eb) a population of black hole binaries of various sizes, having (e) a range of distances from Earth, and created through (e&eb) a variety of scenarios. |
| (262) | They then calculated the amount of gravitational radiation that would result from (e) the merging of these objects. |
| (263) | The team then adjusted this estimate by using (e) details of the GW150914 event—especially the masses of the black holes involved—to make (eb) a most likely guess at the actual distribution of black hole binaries. |
| (264) | The idea is that the observed merger is (=) most likely representative of (e) the general population, rather than being extremely unusual. |
| (265) | If the black hole binary population is represented as (=) a bell curve, or Gaussian, the event should lie within (eb) the main bulge, not in the tails. |
| (266) | The observed mass of GW150914 suggests (eb) that the main population of black hole binaries could be (=) roughly 20 times more massive than expected, which would generate (eb) a background roughly 10 times "louder" than (e) previous estimates. |
| (267) | However, Christensen emphasizes that the result has (e) lots of uncertainty, as it's based on (e) only one event. |
| (268) | "This estimate seems very plausible to me," says astrophysicist Maura McLaughlin of the University of West Virginia in Morgantown. She is a member of the NANOGRAV collaboration, which tries to detect gravitational waves at much lower frequencies than LIGO can observe. "I can't see any holes in their arguments. Unless LIGO got really lucky, then there must be many more of these systems than previously thought." This level of gravitational background implies (eb) that the LIGO and Virgo detectors might be sensitive enough to see it in 2–5 years, sooner than physicists previously expected, says Christensen, which has led to much anticipation. "Everyone is super excited," he says. |
| (269) | Detecting the background would provide (eb) a confirmation of the presence of gravitational waves carrying (e&eb) energy throughout the Universe. It would also be useful, the researchers suggest, in gaining (e) further understanding of precisely how black holes form. |
| (270) | And subtracting this background might someday allow (eb) detection of the fainter waves coming from (e) the expansion of the Universe moments after the big bang. This research is published in Physical Review Letters. –Mark Buchanan Mark Buchanan is a freelance science writer in Normandy, France. |

**GW150914: Implications for the Stochastic Gravitational-Wave Background from Binary Black Holes B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 131102 – Published 31 March 2016**

(271) The LIGO detection of the gravitational wave transient GW150914, from (e) the inspiral and merger of (e) two black holes with masses ≳30M⊙, suggests (eb) a population of binary black holes with (e&eb) relatively high mass.

(272) This observation implies (eb) that the stochastic gravitational-wave background from binary black holes, created from (e) the incoherent superposition of (e) all the merging binaries in the Universe, could be (=) higher than previously expected.

(273) Using the properties of GW150914, authors estimate the energy density of (e) such a background from (e) binary black holes.

(274) In the most sensitive part of the Advanced LIGO and Advanced Virgo band for (e) stochastic backgrounds (near 25 Hz), authors predict $\Omega_{GW}(f=25\ Hz)=1.1^{+2.7}_{-0.9}\times10^{-9}$ with 90% confidence.

(275) This prediction is robustly demonstrated for (e) a variety of formation scenarios with (e&eb) different parameters.

(276) The differences between models are (=) small compared to the statistical uncertainty arising from (e) the currently poorly constrained (e) local coalescence rate.

(277) Authors conclude that this background is potentially measurable by (e) the Advanced LIGO and Advanced Virgo detectors operating at their projected final sensitivity. DOI:http://dx.doi.org/10.1103/PhysRevLett.116.131102 © 2016 American Physical Society

**Study opens new realms of light-matter interaction July 15, 2016 by David L. Chandler. Some sentences might have been omitted due to spatial constraints. Please bear with me.**

(278) Emission spectra are a widely used method for (e) identifying chemical compounds; the bright lines reveal (eb) the different frequencies of light that can be emitted by (e) an atom.

(279) Here, a normal emission spectrum for (e) an atom in a high-energy state (top) is compared to (e&eb) the emission from the same atom placed just a few nanometers (billionths of a meter) away from (e) graphene that has been doped with (e&eb) charge carriers (bottom).

(280) For each energy-level transition, an orange line (or purple cloud) appears if (e) that transition is estimated to be faster than (e) one per microsecond — making it frequent enough to be (=) observed. **Credit: Massachusetts Institute of Technology**

(281) A new MIT study could open up new areas of technology based on (e) types of light emission that had been thought to be (=) "forbidden," or at least so unlikely as to be (=) practically unattainable.

(282) The new approach, the researchers say, could cause (eb) certain kinds of interactions between light and (e&eb) matter, which would normally take billions of years to happen, to take place instead (e&eb) within billionths of a second, under (e&eb) certain special conditions. The findings, based on a theoretical analysis, are reported today in the journal Science in a paper by MIT doctoral student Nicholas Rivera, Department of Physics Professor Marin Soljačić, Francis Wright Davis Professor of Physics John Joannopoulos, and postdocs Ido Kaminer and Bo Zhen.

(283) Interactions between light and matter, described by (e) the laws of quantum electrodynamics, are the basis of a (e) wide range of technologies, including (e) lasers, LEDs, and atomic clocks.

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| (284) | But from a theoretical standpoint, "Most light-matter interaction processes are 'forbidden' by (e) electronic selection rules, which limits (e) the number of transitions between energy levels we have access to," Soljačić explains. |
| (285) | For example, spectrograms, which are used to (e) analyze the elemental composition of materials, show (eb) a few bright lines against (e&eb) a mostly dark background. |
| (286) | The bright lines represent (eb) the specific "allowed" energy level transitions in (eb) the atoms of that element that can be accompanied by (e) the release of a photon (a particle of light). |
| (287) | In the dark regions, which make up most of the spectrum, emission at those energy levels is (=) "forbidden." |
| (288) | With this new study, Kaminer says, "we demonstrate theoretically that these constraints can be lifted" using (e) confined waves within (e&eb) atomically thin, 2-D materials. |
| (289) | "We show that some of the transitions which normally take the age of the universe to happen could be made to happen within (eb) nanoseconds. |
| (290) | Because of this, many of the dark regions of a spectrogram become (=) bright once an atom is (=) placed near a 2-D material." |
| (291) | Electrons in an atom have (e) discrete energy levels, and when they hop from one level to another they give off (eb) a photon of light, a process called spontaneous emission. |
| (292) | But the atom itself is much smaller than the wavelength of the light that gets emitted—about 1/1,000 to 1/10,000 as big—substantially impairing (e&eb) the interactions between the two. |
| (293) | The trick is, in effect, to "shrink" (e) the light so it better matches the scale of the atom, as (=) the researchers show in their study. |
| (294) | The key to enabling a whole range of interactions, specifically transitions in (eb) atomic states that relate to absorbing or emitting light, is (=) the use of a two-dimensional material called graphene, in which light can interact with (e&eb) matter in the form of plasmons, a type of electromagnetic oscillation in the material. |
| (295) | These plasmons, which resemble photons but have wavelengths hundreds of times shorter, are (=) very narrowly confined in the graphene, in a way that makes some kinds of interactions with (e&eb) that matter many orders of magnitude more likely than (e) they would be in ordinary materials. |
| (296) | This enables a variety of phenomena normally considered unattainable, such as (e) the simultaneous emission of multiple plasmons, or two-step light-emitting transitions between (e&eb) energy levels, the team says. |
| (297) | This method can enable the simultaneous emission of (e) two photons that are "entangled," meaning (eb) they share the same quantum state even when (e) separated. |
| (298) | Such generation of entangled photons is (=) an important element in quantum devices, such as those that might be used for cryptography. |
| (299) | Making use of these forbidden transitions could open up (eb) the ability to tailor the optical properties of materials in ways that had not been thought possible, Rivera says. |
| (300) | "By altering these rules" about the relationship between light and (e&eb) matter, "it can open (eb) new doors to reshaping (e&eb) the optical properties of materials." |
| (301) | Kaminer predicts that this work "will serve as (=) a founding piece for (e) the next generation of studies on light-matter interactions" and could lead to (eb) "further theoretical and experimental advances in many fields which rely on (e) light-matter interactions, including (e) atomic, molecular and optical physics, photonics, chemistry, optoelectronics, and many others." |

(302) Beyond its scientific implications, he says, "this study has possible applications across multiple disciplines, since (e) in principle it has potential to enable the full use of (e) the periodic table for optical applications."

(303) This could potentially lead to (eb) applications in spectroscopy and sensing devices, ultrathin solar cells, and new kinds of materials to absorb solar energy, organic LEDs with higher efficiencies, and photon sources for (e) possible quantum computing devices."From the standpoint of fundamental science, this work lays the groundwork for a subfield that just a few years ago was difficult to imagine and until now was largely unexplored," Soljačić says.

(304) "Two-dimensional materials confine (e) fields to a surface and motion to a plane, making possible (eb) many effects that are orders-of-magnitude too weak to (e) appear in a bulk volume," says Jason Fleischer, an associate professor of electrical engineering at Princeton University, who was not involved in this research.

(305) This work, he says, "systematically explores how 2-D materials improve light-matter interactions, laying a theoretical foundation for faster electronic transitions, enhanced sensing, and better emission, including the compact generation of broadband and quantum light." More information: N. Rivera et al: Shrinking light to allow (eb) forbidden transitions on the atomic scale, Science (2016). DOI: 10.1126/science.aaf6308 Journal reference: Science search and more info website: Provided by: Massachusetts Institute of Technology search and more info website. This story is republished courtesy of MIT News (web.mit.edu/news office/), a popular site that covers news about MIT research, innovation and teaching.

**Viewpoint: Peering Back to the Beginning of Time Lawrence M. Krauss, Department of Physics and School of Earth and Space Exploration, Arizona State University, Tempe, AZ 85287, USA and Mount Stromlo Observatory, Research School of Astronomy and Astrophysics, Australian National University, Weston, ACT 2611, Australia June 19, 2014• Physics 7, 64**

(306) The BICEP2 collaboration reports the detection of BB-mode polarization of (e) the cosmic microwave background—a signal that might originate from (e) gravitational waves created by (e) inflation during the very earliest moments in the evolution of our Universe. Figure caption expand figure P. A. R. Ade et al. [1]; Image on homepage: BICEP2 Collaboration (deleted due to restrictions on space)

(307) Figure 1: (Top) BB-mode polarization pattern observed by (e) the BICEP2 telescope [1].

(308) The twisting pattern in the polarization of (e) the cosmic microwave background (CMB) might originate from (e) gravitational waves created by inflation in the very first moments of the universe.

(309) Inflation—the hypothesis that the Universe underwent (e&eb) a phase of superluminal expansion in a brief period following the big bang—has (e) the potential of explaining, from first principles, why the Universe has (e) the structure we see today.

(310) It could also solve outstanding puzzles of standard big-bang cosmology, such as why the Universe is, to a very good approximation, (=) flat and isotropic (i.e., it looks the same in all directions).

(311) Yet we do not yet have a compelling model, based on (e) fundamental particle physics principles, that explains inflation.

(312) And despite its explanatory power and a great deal of suggestive evidence, we still lack an unambiguous and direct probe of (e) inflation.

(313) Theorists have developed different models for inflation, which all share (e&eb) a common, robust prediction: Inflation would have created (eb) a background of gravitational waves that could have (eb) an observable effect.

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| (314) These waves would cause (eb) subtle, characteristic distortions of (e) the cosmic microwave background (CMB)—the oldest light in the Universe, released when (e) photons decoupled from matter and the Universe became (=) transparent to radiation |
| (315) Until now, such a background of gravitational waves has (e) escaped detection. It is therefore very exciting that the BICEP2 (Background Imaging of Cosmic Extragalactic Polarization) collaboration has now reported in Physical Review Letters the observation of polarization modes of CMB radiation [1] that are consistent with (e&eb) such a background. |
| (316) If confirmed, BICEP2's finding would provide (eb) the first direct probe of physical processes that occurred a mere 10−3610-36 seconds after (e) the big bang. |
| (317) The result also offers a rare opportunity to directly test theoretical ideas, including inflation and grand unification (the theory that merges the electromagnetic, weak, and strong forces into a single force), at energies 13 orders of magnitude greater than (e) we can currently access at the Large Hadron Collider. |
| (318) Inflation is based on (e) the possibility that a scalar field in the early Universe got momentarily stuck in (eb) a "false vacuum state"—a metastable condition in which the Universe wasn't at (eb) its global potential minimum. |
| (319) The energy stored in this vacuum state could drive (eb) an exponential expansion of spacetime (similar to, but much faster than, the expansion due to dark energy seen today) before (e) a phase transition to the true vacuum state occurs, releasing (eb) this energy into (e&eb) particles and radiation. |
| (320) This creates (eb) an initially hot state for (e) a subsequent big-bang-like expansion from (e) which our Universe subsequently evolved [2]. |
| (321) If the Universe grew by at least a factor of 10261026, (eb) inflation can explain (eb) both the apparent flatness and isotropy of the Universe [2]: |
| (322) The expansion would stretch (eb+ eb-dot) any initial curvature of (e) the Universe to near flatness, and smooth out (e&eb) any initial inhomogeneities, just as stretching (eb+ ebdot) a piece of fabric would do. |
| (323) Isotropy is justified by the fact that the Universe, before inflation, was once much smaller and could have been contained in (eb) a causally connected region that was in thermal equilibrium, resulting in (eb) the approximately uniform, isotropic temperature distribution observed on (eb) large scales today. |
| (324) Importantly, the theory can explain the (eb) origin of structure in the Universe: Microscopic quantum fluctuations associated with (e&eb) the scalar field driving inflation could have been "frozen in" during the exponential expansion, resulting, (eb) after inflation, in density fluctuations that might have seeded all observed large-scale structures in the Universe, including (e) all visible galaxies and clusters of galaxies. |
| (325) The signs of such primordial fluctuations have indeed been detected through (e&eb) their effect on (e&eb) the CMB, first by the Cosmic Background Explorer (COBE) [3] and later by other missions. |
| (326) These observational studies detected small temperature variations (anisotropies) in (eb) the CMB and found a spectrum compatible with (e&eb) that predicted by inflation. |
| (327) The magnitude of these fluctuations suggests (eb) an energy scale of inflation (loosely speaking, the energy stored throughout space in (eb) the scalar field potential that drove inflation) close to (e) 10161016 giga-electron-volts (GeV). |
| (328) But while consistent with inflation, such fluctuations are not (e) conclusive evidence: Their predicted shape and magnitude is known to vary with (e&eb) the details of the inflationary potential (whose form is not yet experimentally constrained). |

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| (329) Hence many different observations could be accommodated, and any agreement between theory and observation is at best suggestive. There is also a possibility that other, noninflationary sources might contribute. However, gravitational waves—the propagating ripples in spacetime predicted by Einstein in 1916—might provide (eb) direct evidence of inflation. |
| (330) During inflation, gravitational fields would also fluctuate because of (e) quantum effects. |
| (331) These fluctuations would similarly get "frozen in" and re-emerge after inflation as (=) classical gravitational waves, whose amplitude depends only on (=) the energy stored in the field driving (eb) inflation, and not on other model-dependent details [4]. |
| (332) If the scale of inflation were indeed as high as a few times 10161016 GeV, this background could produce (eb) specific temperature anisotropies in the CMB [5]. |
| (333) These anisotropies have thus far not been detected by (e) the most recent searches by the WMAP and PLANCK satellites, and their absence was used to (e) put upper limits on the scale of inflation somewhat below 10161016 GeV. |
| (334) But a more sensitive CMB probe of (e) gravitational waves exists. |
| (335) Large-wavelength gravitational waves produced by (e) inflation would have alternatively compressed and expanded (e&eb) space, producing (eb) spatial variations in (eb) the intensity of radiation bombarding free electrons in (eb)the early Universe (see, e.g., Ref: [6]). |
| (336) As photons decoupled from (e) matter and CMB radiation was released, scattering from (e) the undulating sea of electrons would have created (eb) patterns in the polarization of CMB. |
| (337) Such polarization anisotropies are predicted in (eb) inflationary models to be (=) most visible when (e) looking at the sky on spatial scales comparable to (e&eb) or somewhat larger than the size of (e) the Universe's horizon at (eb) the time when the CMB was created. |
| (338) This corresponds to (e&eb) about 1 degree of the field of view of the sky. |
| (339) Over the past decade, various CMB polarization experiments have thus chased the observation of polarization anisotropies on (e&eb) these scales but, so far, also yielded (eb) merely upper limits on a possible gravitational background. |
| (340) All of this has now changed with (e&eb) the BICEP2 collaboration's [1] report of (e) a signal that has the predicted (eb) characteristics of such a background at precisely the angular scales where (e) the inflationary signal is expected (see Fig. 1, top deleted due to spatial restraints). |
| (341) The BICEP2 Collaboration uses (e) a microwave telescope located at the South Pole, with (e&eb) an array of polarization-sensitive detectors pointing at (eb) a small field (a solid angle of about 11 percent of the full sky) with a resolution of approximately 0.50.5 degrees, thus sufficient to probe the relevant angular scales for the detection of (e) polarization anisotropies originating from (e) **inflation**. |
| (342) BICEP2's search is based on (e) two expected features of primordial gravitational waves. |
| (343) First, gravitational waves would produce (eb) a characteristic spatial polarization pattern on CMB—a consequence of (e) their quadrupole nature. |
| (344) One can decompose (e&eb) any spatial polarization pattern into (e&eb) two orthogonal components, a "curl-free" component and (e&eb) a component involving a twisting pattern. |
| (345) In analogy to electromagnetic modes, these polarization patterns are called "EE modes" and (e&eb) "BB modes," respectively (see Fig. 1). |
| (346) Scalar fluctuations in the density of (e) matter would produce (eb) EE modes, while (e) gravitational waves would instead generate (eb) a BB-mode pattern. |

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| (347) The second feature is the aforementioned characteristic spatial scale of (e) expected CMB polarization signatures, including (e) wavelength scales comparable to (e&eb) or larger than the horizon scale at the time the CMB was created. |
| (348) Since inflation would produce (eb) both gravitational waves and (e&eb) scalar density perturbations, BICEP2, like all CMB-polarization experiments, quote their results in terms of (e&eb) the ratio rr of the amplitude of (e) tensor perturbations (associated with (e&eb) gravitational waves) to (e) the amplitude associated with (e&eb) scalar density fluctuations (reflected in (eb) the magnitude of CMB temperature fluctuations measured by (e) experiments like COBE). |
| (349) A nonzero rr implies (eb) the detection of a BB-mode pattern and primarily measures (eb) the strength of any possible gravitational wave signal which, if (e) it comes from inflation, is then determined by (e) its energy scale (e.g., see Ref. [7] for a review). |
| (350) BICEP2 reported a nonzero value of r=0.2(+0.07−0.05)r=0.2(+0.07-0.05), on scales ranging from (e) the horizon scale up to 4–5 times larger, with (e&eb) a claimed significance of better than 5σ5σ. |
| (351) If confirmed, a value of 0.2 corresponds to (e&eb) an energy density during (e&eb) inflation of ∼1.2×1016∼1.2×1016 GeV. |
| (352) The claimed signal is far larger than expected, partly because (e) the PLANCK satellite, despite not having yet reported its polarization measurements, has (e) previously put an upper limit on (eb) rr (<0.11<0.11). PLANCK's estimate is based on (e) the absence of a gravitational wave signature on (eb) the CMB temperature anisotropy on (eb) its largest scales (PLANCK observes the whole sky and therefore angular scales up to 360360 degrees). |
| (353) The apparent discrepancy between the PLANCK upper bound and (e&eb) the BICEP2 detection might be a cause for concern. |
| (354) However, the two bounds probe different spatial scales, and the BICEP collaboration has noted that, while tensor perturbations are expected to be (=) roughly scale independent, some scale dependence of (e) scalar density perturbations could relax (e&eb) the PLANCK constraint and bring (eb) the two results into agreement. |
| (355) Of greater concern, of course, is the possibility that the signal could be determined by (e) foreground contamination—any source along (e&eb) the line of sight of the detector that could distort or contaminate (e&eb) the primordial CMB signal. |
| (356) BICEP2 tested for (e) a wide variety of possible foregrounds, claiming (eb) they cannot be responsible for (e) the observed signal. |
| (357) One obvious candidate arises from (e) gravitational lensing: The intervening mass of (e) astrophysical bodies could distort (e&eb) the spatial pattern of polarization in a way that that could transform (e&eb) an EE-mode signal into (e&eb) a BB-mode signal. |
| (358) The BICEP collaboration has analyzed this possibility extensively. Based on (e) measurements of (e) CMB temperature (and hence density perturbations) and on (e&eb) models of the mass distribution, they argue that such effects would only be (=) significant on smaller angular scales than those contributing to their measurements of rr. |
| (359) But since BICEP2's first announcement on March 17, 2014, a number of scientists have raised another foreground concern due to (e) the possible impact of (e&eb) our Galaxy's dust. |
| (360) BICEP2 modeled dust contamination scenarios based on (e) a variety of theoretical models, using data from (e) its predecessor BICEP1 and preliminary dust estimates from (e) the PLANCK satellite. |
| (361) They concluded that dust could not (e) reproduce the magnitude of the observed signal. |

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| (362) However, PLANCK has in the meanwhile released (eb) more complete measurements of variation in polarized dust foregrounds in (eb) many regions of the Galaxy. |
| (363) Their analysis increases the uncertainty associated with (e&eb) the dust contribution to (e) the BICEP2 result and also accommodates (e) a potentially larger polarized dust contribution. |
| (364) Some researchers argued that dust polarization at a level 2–3 times higher than that assumed by BICEP2 is possible and could therefore plausibly be responsible for (e) the entire signal (e.g., see Ref: [8]). |
| (365) While the BICEP2 collaboration continues to maintain that this is extremely unlikely based on (e) their estimates, many other researchers are more skeptical. |
| (366) In any case, a variety of upcoming measurements, from PLANCK and other detectors are (=) expected to resolve this important ambiguity within the coming year. |
| (367) Even if the BICEP2 signal does arise from (e) primordial gravitational waves, various consistency checks will be required (e) before it can be definitively argued that the signal arises from (e) inflation. |
| (368) Inflation is certainly the best-motivated possibility, and the spectral shape of the signal, at least for (e) the five largest-scale modes observed by BICEP2, agrees extremely well with (=) inflationary predictions (see Fig. 1, bottom). |
| (369) Several results have already appeared that strongly constrain (e) other possible sources, ruling out (e) for example, the possibility of gravitational waves generated by (e) other possible phase transitions after inflation [9] |
| (370) . This conclusion is reinforced by more careful analyses of the scale dependence of (e) the predicted signal (e.g., Ref. [10]). |
| (371) But it will nevertheless be important to assure that all possible alternative sources of (e) gravitational waves are ruled out. |
| (372) If it can be established that the BICEP2 represents (eb) a gravitational wave signal from (e) inflation, there is one additional exciting implication. |
| (373) Standard calculations suggest (eb) that the gravitational wave signal from inflation is due to (e) quantum fluctuations in the gravitational field, and this appears to require (e) the existence of gravitons—the hypothetical elementary particles that mediate (e&eb) the force of gravity. |
| (374) Research by Frank Wilczek and myself, based on (e) a dimensional analysis argument [11], suggests that,(eb) independently of the method used to calculate the gravitational wave spectrum, a quantum gravitational origin is (=) required. |
| (375) If so, the BICEP2 result then implies (eb) that gravity is ultimately a quantum theory, a result of fundamental importance for physics. |
| (376) It is hard to imagine a more important or wider-ranging cosmological observation. If confirmed by upcoming experiments, the BICEP2 results will push the frontiers of empirical physics forward in an unprecedented way, offering (eb) the first direct observation of gravitational waves and of quantum gravity effects as well as a new probe of the earliest moments in the Universe's life and of physics at grand-unified energy scales. This will thus constrain both our understanding of fundamental interactions and the evolution of our Universe as a whole. |

**Detection of B-Mode Polarization at Degree Angular Scales by BICEP2 P. A. R. Ade et al. (BICEP2 Collaboration) Phys. Rev. Lett 112, 241101 – Published 19 June 2014**

(377) Authors report results from the BICEP2 experiment, a cosmic microwave background (CMB) polarimeter specifically designed to search for (e) the signal of inflationary gravitational waves in (eb) the B-mode power spectrum around $\ell \sim 80$.

(378) The telescope comprised a 26 cm aperture all-cold refracting optical system equipped with (e&eb) a focal plane of 512 antenna coupled transition edge sensor 150 GHz bolometers each with temperature sensitivity of $\approx 300\ \mu K_{CMB}\sqrt{s}$. BICEP2 observed from (e) the South Pole for three seasons from 2010 to 2012.

(379) A low-foreground region of sky with (e&eb) an effective area of 380 square deg was observed to (e) a depth of 87 nK deg in Stokes Q and U.

(380) In this paper authors describe the observations, data reduction, maps, simulations, and results. They find an excess of B-mode power over the base lensed-ΛCDM expectation in the range $30<\ell<150$, inconsistent with (e&eb) the null hypothesis at (eb) a significance of $>5\sigma$.

(381) Through jackknife tests and simulations based on (e) detailed calibration measurements authors show that systematic contamination is (=) much smaller than the observed excess.

(382) Cross correlating against WMAP 23 GHz maps we find that Galactic synchrotron makes a negligible contribution to the observed signal. We also examine a number of available models of polarized dust emission and find that at their default parameter values they predict power $\sim(5–10)\times$ smaller than the observed excess signal (with no significant cross-correlation with our maps). However, these models are not sufficiently constrained by external public data to exclude the possibility of dust emission bright enough to explain the entire excess signal. Cross correlating BICEP2 against 100 GHz maps from the BICEP1 experiment, the excess signal is confirmed with $3\sigma$ significance and its spectral index is found to be consistent with that of the CMB, disfavoring dust at $1.7\sigma$. The observed B-mode power spectrum is well fit by a lensed-ΛCDM+tensor theoretical model with tensor-to-scalar ratio $r=0.20+0.07-0.05$, with $r=0$ disfavored at $7.0\sigma$. Accounting for the contribution of foreground, dust will shift this value downward by an amount which will be better constrained with upcoming data sets. Received 4 April 2014 DOI:http://dx.doi.org/10.1103/PhysRevLett.112.241101 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**(383) Physicist claims to have observed quantum effects of Hawking radiation in the lab for the first time August 16, 2016 by Bob Yirka report**

(384) This computer-simulated image shows a supermassive black hole at the core of a galaxy. The black region in the center represents the black hole's event horizon, where no light can escape the massive object's gravitational grip. The black …more

(385) (Phys.org)—Jeff Steinhauer, a physicist at the Israel Institute of Technology, has published a paper in the journal Nature Physics describing experiments in which he attempted to create a virtual black hole in the lab in order to prove that Stephen Hawking's theory of radiation emanating from black holes is correct —though his experiments are based on sound, rather than light. In his paper, he claims to have observed the quantum effects of Hawking radiation in his lab as part of a virtual black hole—which, if proven to be true, will be the first time it has ever been achieved.

(386) For many years, scientists believed that nothing could ever escape from a black hole. But in 1974, Stephen Hawking published a paper suggesting that something could—particles that are now called Hawking radiation. His idea was that if a particle (and its antimatter mate) appeared spontaneously at the edge of a black hole, one of the pair might be pulled into the black hole while the other escaped, taking some of the energy from the black hole with it—which would explain why black holes grow smaller and

eventually disappear. Because such emissions are so feeble, no one has been able to measure Hawking radiation, so researchers have instead tried to build virtual black holes in labs to test the theory. One type of virtual black hole was proposed back in 1981 by Bill Unruh with the University of British Columbia—he suggested that an analogue might be created using water instead of light. He imagined a phonon existing at the edge of a waterfall—as the water speeds up, it begins to move faster than the speed of sound, causing it to be trapped. But if the phonon had an entangled mate that eluded the fall by moving away before getting caught up, it could escape. In this new effort, Steinhauer has built a device based on that idea and in so doing, claims he has observed an analogue of Hawking radiation.

(387) The experiment consisted of creating an entangled pair of phonons sitting inside a bit of liquid that had been forced (via laser) to move very fast and then observing the action as one of the pair was pulled away as the liquid began to move faster than the speed of sound, while the other escaped—the fluid was a Bose-Einstein condensate of rubidium-87 atoms. After repeating the experiment 4,600 times Steinhauer became convinced that the particles were entangled, a necessity for a Hawking radiation analogue. His findings do not prove Hawking's theory to be true, of course, but they do appear to add a degree of credence that other researchers have thus far not been able to achieve. More information: Jeff Steinhauer. Observation of quantum Hawking radiation and its entanglement in an analogue black hole, Nature Physics (2016). DOI: 10.1038/nphys3863

**(388) Abstract**

We observe spontaneous Hawking radiation, stimulated by quantum vacuum fluctuations, emanating from an analogue black hole in an atomic Bose–Einstein condensate. Correlations are observed between the Hawking particles outside the black hole and the partner particles inside. These correlations indicate an approximately thermal distribution of Hawking radiation. We find that the high-energy pairs are entangled, while the low-energy pairs are not, within the reasonable assumption that excitations with different frequencies are not correlated. The entanglement verifies the quantum nature of the Hawking radiation. The results are consistent with a driven oscillation experiment and a numerical simulation. Journal reference: Nature Physics search and more info website Read more at: http://phys.org/news/2016-08-physicist-quantum-effects-hawking-lab.html#jCp

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending

upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

## DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t-T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

These equations can be written for

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(1) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.

(2) However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].

(3) Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.

Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time

delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

**TIME DELAY**

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types system takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t-T_7)F_7(R_7(t-T_7) - N_7(t)M_7((R_7(t) \qquad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7} + \frac{N_7 dG_7}{dR_7} - \frac{dB_7}{dR_7} > \frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7} + \frac{dM_7}{dR_7}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also normalized by $M$, which encapsulate effects acting on the credit condition, one derived from

the time-delayed or "legged" processes ($L_7$) and the other derived from the non-delayed or "immediate" processes($I_7$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_7 = \frac{N_7 G_7}{2M_7} \frac{d(7nF_7)}{dR_7} (\geq 0) \quad (7a)$$

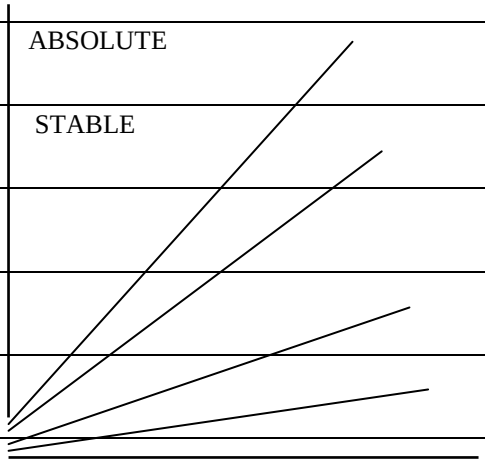
$$I_7 = -\frac{N_7 G_7}{2M_7} \frac{d(7nM_7}{dR_7} (\geq 0) \quad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7 > L_7 - I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_7 > L_7 - I_7$, when $P_7^2/4 > I$ (8a)

$P_7 > L_7 - I_7 + \frac{(I_7 - P_7^2}{(L_7 - I)}$ when $P_7^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,



**DELAYED PREDATOR CREDIT TYPES EFFECT WITH ITS WIDE RANGING RAMIFICATION**

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

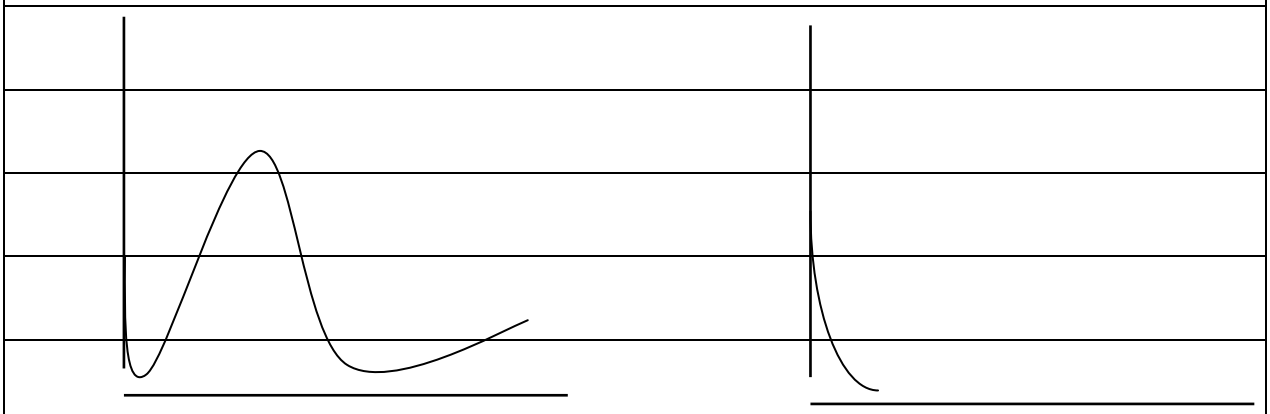
It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always

decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



**TIME DELAY**

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

To scores of problems this relationship is important. Of them some of the examples are as follows:

**New look at the galaxy-shaping power of black holes**

(1) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving.

(2) "These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science.

(3) "Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy."

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| (4) | This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return. |
| (5) | The remaining matter gets absorbed into the black hole, adding to (e) its mass. |
| (6) | The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne. |
| (7) | The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution. |
| (8) | Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form. |
| (9) | "Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara. |
| (10) | "It creates (eb) these bubbles, and the bubbles keep the plasma hot. |
| (11) | That's what prevents galaxies from (e) becoming even bigger than they are now." |
| (12) | Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space. |
| (13) | In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene. |
| (14) | Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve. |
| (15) | "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara. |
| (16) | "It's like nature's thermostat, which keeps these galaxies from (e) growing. |
| (17) | If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars." |

**A new look at the galaxy-shaping power of black holes**

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| (18) | An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the |
| (19) | McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said. |
| (20) | Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure. |

**Doomed satellite missions**

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| (21) | Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving. |
| (22) | At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon. |

(23) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.

(24) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.

(25) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.

(26) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people

(27) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."

(28) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.

(29) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(30) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.

(31) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.

(32) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.

(33) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.

(34) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well.

(35) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma.

(36) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6.

(37) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of $164 \pm 10$ kilometres per second in the region 30–60 kiloparsecs from the central nucleus.

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| (38) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (39) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (40) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (41) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (42) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (43) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (44) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (45) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (46) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (47) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (48) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (49) We know dark matter exists only from (e) its gravitational effects. |
| (50) It holds together galaxies and (e&eb) galaxy clusters. |
| (51) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (52) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |
| (53) They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (54) X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (55) They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (56) This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (57) An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |

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| (58) Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (59) "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (60) The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (61) Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (62) Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (63) "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (64) Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (65) "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (66) Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories. |
| (67) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen. |
| (68) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer. |

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| (69) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants. |
| (70) These were treated as (=) "aware" and "unaware" trials, respectively. |
| (71) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it. |
| (72) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another. |
| (73) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion. |
| (74) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness. |
| (75) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said. |
| (76) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness." |
| (77) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner. |
| (78) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain. |
| (79) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said. |
| (80) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

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| (81) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter. |
| (82) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined. |
| (83) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions. |

(84) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.

(85) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.

(86) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.

(87) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.

(88) Noda and his colleagues instead studied the change in emissions over (e&eb) time.

(89) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.

(90) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).

(91) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.

(92) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.

(93) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(94) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking "palm tree" organization.

(95) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society

(96) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible?

(97) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem.

(98) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other.

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| (99) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible. |
| (100) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges. |
| (101) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century. |
| (102) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching. |
| (103) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom. |
| (104) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth. |

**Spheres on trees**

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| (105) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration. |
| (106) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement. |
| (107) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum. |
| (108) By visualizing the problem in this way, the researchers noticed that these particular graphs don't (e) have lots of branches extending from lots of (e) other branches. |
| (109) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure. |
| (110) The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network. |
| (111) As a result, it provides (eb) an important clue for finding the global minimum. |
| (112) It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum. |
| (113) This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large. |

**Thompson's problem 2**

(114) By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society

(115) The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima.

**(116)** The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node.

**Implications for Thomson's and Smale's problems**

(117) Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before.

(118) "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around."

(119) As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem.

(120) "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view."

(121) In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges.

(122) From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear.

(123) "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK.

(124) "For larger numbers, there will likely be (=) multiple funnels.

(125) Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum."

(126) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.

(127) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.

(128) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t-T_7)F_7\big(R_7(t-T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t-T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) and the effective prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Look at the resemblance of String Theory and Abhinava's earlier theory and how the model could be applied: The world of manifestation and differentiated perceptions (Yikalpa) thus extends from one Centre to (eb) the next. **What is meant again is the realisation or the evolution of individual consciousness or general ledger to be (=)the cosmic general ledger is the sole criterion for the Satchidananda state where all dualities cease** Although it is never in fact divorced from (e) the subject who resides there, the ignorant fail (e) to grasp this fact and so, cut off from (e) the Centre, the world of objectivity becomes (=) for them the sphere of Maya. Bhagavatotpala quotes the Light of Consciousness amvitprakdsa): This ever pure experience (suddhanubhava) is variegated by (e) each form [revealed within it] Even so it remains unstained (nirmala) when moving to another. Just as a cloth which is naturally white, once dyed, cannot (e) change colour without [first] becoming white again, similarly the pure power of awareness, (citi) once coloured by (e) form, is pure [again] at the Centre where that form is (=) abandoned and from whence it proceeds to (e&eb) another. In his Essence of Vibration (Spandasarpdoha), K§emaraja explains that the rise and fall of every individual perception in the field of awareness is (=) a specific pulsation of consciousness. This is exactly the individual general ledger transaction lifted to the transaction in cosmic general ledger. During The Initial Instant Of Perception, T Consciousness Is Manifestly Apparent And The Yogi, Participating In Its Plenitude, Observes The Outer World Without Being Attached To Any Particular Divine Body And Sacred Circle Of The Senses Singling It Out From Any Other, Like A Man Who Observes A City From A High Mountain Peak. He Sees The Outer World Reflected Within His Consciousness Free Of Thought-Constructs And So 'Stamps* The Outer On The Inner While Absorbing The Object And Means Of Knowledge In The Pure Subject Which Grasps Them As The Expansion Of His Own Nature. K§emaraja Says: By Penetrating Into Bhairavimudra, The Yogi Observes The Vast Totality Of Beings Rising From, And Dissolving Into, The Sky [Of Conscious- Ness], Like A Series Of Reflections Appearing And Disappearing Inside A Mirror. Through (e&eb) The Practice Of Bhairavimudra The Yogi Realises That He Is (=) The Substratum Consciousness (Adhisfhatr) Which Both Underlies

(e&eb) And Is (=) The Essence Of All Things. He Discovers That Phenomena Have (e) No Independent Existence Apart From (e) Him And So Are, In This Sense, Void. At The Same Time, He Realises That Because (e) All Things Are Consciousness, They Are (=) Far From Unreal. He Views The Outer World Yet Sees (e) It Not. Beyond Both Voidness And Non-Voidness He Penetrates Into (e&eb) The Supreme Abode (Pararn Padam) Of Siva's Consciousness. [The Powers Of The Senses] Endowed With (e&eb) The Attributes Of The Great Union [Between Subject And (e&eb) Object] Whose Form Is The Wakening Of Man's Spiritual Potential (Kuntfalinf), Fill (e&eb) [With Consciousness] The Outer Clatter Of Diversity (Bhedatfambara) Born Of (e) Its Intense Power And Are Then Established In (eb) The Unobscured Abode Of The Void Of Consciousness To Shine [There] Eternally. Thus Residing Beyond (e) Being And Non-Being, The Sole Protector Of The Unity Which Is Tranquil And Expanding [Consciousness], Whose Glory Is All-Embracing And Form Unobscured, Is Called Bhairavimudra. Through (e&eb) The Practice Of Bhairavimudra, The Yogi Unites The Universal Vibration Of T Consciousness With (e&eb) The Individual Pulsation Of Objectivised 'This' Consciousness. The Two Aspects Of Consciousness Are Now In (eb) A State Of Equilibrium Like (=) The Two Pans Of An Evenly Weighed Balance And The Yogi Experiences (e&eb) The Pure Knowledge (Suddhavidya) That: 'I Alone Am (=)All Things'. Thus Becoming (=) The Master Of The Wheel Of Energies He Is (=) Free, Like Siva, To Create And Destroy. When [The Yogi] Is Well Established, Without (e) Wavering, Solely In The Integral Egoity Of His Authentic Nature, The Spanda Principle, And Is Absorbed In Contemplation (Samavisfa), He Becomes (=) One With It The Doctrine Of Vibration (Tanmaya). Then . . . Dissolving And Creating The Universe By Means Of (er) His Introverted And Extroverted Absorption, He Destroys And Creates (e&eb) All Things Out Of Sarikara, His Innate Nature. [Thus] He Assumes The State Of The Universal Experiencer And Having Absorbed All That Is To Be Experienced From [The Grossest Level] — Earth — To (The Subtlest) — Siva — He Reaches (eb) The State Of The Supreme Subject By (e) Progressively Recognising [His Identity With Him]. Thus, Introverted And Extroverted Absorption Both Lead To (eb) The

Recognition Of The Pulsation (Spanda) Of (e) One's Own Consciousness. This Is Where (e) Individual General Ledger Becomes (=) Cosmic General Ledger. At The Level Of Consciousness Corresponding To (e&eb) Siva's Basic State (Saryibhavdvastha), The Alternation From Inner To (e&eb) Outer Is Instantaneously Resolved Into (e&eb) The Vibration Of His Nature. When The Yogi Finally Comes To Be (=) Constantly Aware Of This Reality, His Enlightenment Is (=) Full And Perfect. Freed Of All Means (Anupaya) And Delighting In The Power Of His Bliss (Anandasakti), He Knows And Does Whatever He Pleases. The Yogi Seeking Self-Realisation Must Acquire (e) Mastery Over This Movement. Ksemaraja Stresses That The Doctrine Of Vibration Teaches That Liberation Can Only Be Achieved By (e) First Withdrawing All Sense Activity In Introverted Contemplation (Nimilanasamadhi) To (e) Then Experience The 'Great Expansion' (Mahdvikdsd) Of (e) Consciousness While (e) Recognising This To Be (=) A Spontaneous Process Within (eb) It. This Is Done Through (e&eb) The Practice Of Kramamudra. A Passage From The Now Lost Kramasiitra Explains: Although The Adept's Attention [May Be] Outwardly Directed, (eb) He Enjoys Contemplative Absorption Through (e&eb) The Introverted Aspect Of (e) Kramamudra. Initially He Turns Inward From (e)The Outside World And [Then] From Within [Himself] He Exits Into The Outer World Under The Influence Of His Absorption. There Is No Necessity Of Any Glorification Or Mortification Here. It Is Just Like Joining A Job And Retiring And Realising All That Happened Is Just By His Own Effort And Nothing Else. Thus The Sequence (Krama) In This Attitude (Mudra) [Ranges Through] Both Inner And Outer. The Yogi Must Pervade The Surface Level Of Awareness (Yyutthana) With The Same Bliss He Experiences Plunging Into The Depths Of Con- Templative Absorption (Samadhi). Note Here That Samadhi Is A Death Like State And Still Being Matter And Information Stored You Experience Consciousness. Submerging Himself And Emerging Repeatedly From Samadhi, He Eventually Recognises That The Unity Of Consciousness Pervades Both States: The Best Of Yogi's, Who Has Achieved A State Of Complete Absorption Even When Risen From Meditation, [Inwardly] Vibrating Like A Drunkard In Blissful Inebriation From The After-Effects Of The Nectar Of Con- Templation, Sees All Things

Dissolving In The Sky Of Consciousness Like A Cloud In The Autumn Sky. He Plunges Repeatedly within Himself Divine Body and Sacred Circle of the Senses and Becomes (=) Aware Of His Identity With (e&eb) Consciousness By (e) The Practice Of Introverted Contemplation. Thus Even When He Is Said To Have Risen From (e) Absorption, He Is (=) One With [His] Experience Of It. Ksemaraja Goes On To Explain That This Practice Is Called 'Mudra' Because (e) It Both Fills The Adept With (e&eb) Bliss (Mud) And Is Itself The Bliss Of (e) Consciousness. Moreover, It Dissolves (e) Away (Dra) All Bondage And 'Stamps' The Universe Of Experience With (e&eb)The Seal (Mudra) Of The Fourth State (Turiya) Of Enlightened Consciousness Beyond, And Including, The Three States Of Waking, Dreaming And Deep Sleep. It Is Called 'Krama' Because (e) It Is The Root Source Of (e) All Emanation And All Other Conscious Processes Which Succeed (e&eb) One Another In Ordered Sequence (Krama) And Is, (=)At The Same Time, Their Successive (Krama) Appearance As Well. By The Practice Of Kramamudra The Opposites Fuse And Siva And Sakti Unite. They Yogi Comes To Experience The Simultaneous Pervasion Of All The Lower, Grosser Categories Of Existence By The Higher And The Presence Of The Lower In The Higher. Commencing His Practice In A Low Form Of Bhairavimudra, The Yogi Conjoins The Outer With The Inner; Then, In Kramamudra, He Fills Both The Outer With (e&eb) The Inner And (e&eb) The Inner With (e&eb)The Outer. When He Achieves (eb) Perfection In This Two-Fold Movement, He Attains To (e) The Highest Form Of Bhairavimudra In Which The Two Merge Completely In (eb)The Experience Of (e) The Absolute (Anuttara), Free Of (e) All Differentiation And (e&eb) Polarities. If He Fails To Maintain Awareness Of This State, He Again Falls Into (eb) Kramamudra Until He Has Finally Completely Merged All The Highest States In (e&eb) The Lower And The Lower In (e&eb) The Higher. He Then No Longer Needs To (e) Resort To (e) Any Means (Anupaya) To (e) Achieve Liberation. All He Says or Does Anything He Perceives or Thinks, Instantly Occasions In (eb) Him The Highest Level Of Consciousness. Thus The Fruit Of Bhairavimudra Is (*=)The Wonder (Camatkara) Or Amazement (Vismaya) That Overcomes (e)The Yogi When (e) He Reaches The Plane Of (e) Union (Yogabhumika) Y Xsl

Where (e) All Opposites Merge In (e&eb)The Radiance Of The Great Light Of Consciousness. The Stanzas On Vibration Teach: How Can One Who, As If Astonished, Beholds (e) His Own Nature As (=) That Which Sustains [The Existence Of Everything] Be Subject To (e&eb) This Painful Round Of Transmigration (Kusrti)? The Yogi, Recognising His True Nature To Be (=)The Supreme Subject, Is Astonished To (e) Suddenly Discover That The Individual He Thought He Was, Caught Up (e&eb) In The Trammels Of Thought And Living In A World Enmeshed In (e&eb) The Web Of Time And Space, Does Not Really Exist At All. He Experiences A 'Turning About' (Paravrtti) In (eb)The Deepest Seat Of The Doctrine Of Vibration Consciousness As He Penetrates His True Nature. The Sudden Eruption Of This Intuition (Pratibha) Arouses In Him A Cry Of Amazement As He Transcends All Thought-Constructs And, Perfectly Absorbed In His Own Nature, Is Liberated. The Path To Liberation Essentially, Spanda Doctrine Is Concerned With Two Matters. The First Is To Impart To Those Who Are Fit To Receive The Teachings A Deeper Understanding Of The Ultimate Goal Of Life (Upeya). When We Have Understood What Truly Benefits Us And Is Worth Attaining And What, On The Contrary, Is Of No Real Value But Stands In The Way Of This Attainment, We Can Begin To Make Progress Towards Our Goal. This Is Spanda Doctrine's Second Concern, Namely, To Show The Way In Which We Can Develop Spiritually Through Siva's Grace And The Right Application Of The Means To Realisation That It Teaches. When Both These Aspects Of The Teaching Have Been Correctly Understood And Applied, The Spanda Yogi Achieves A Clear And Permanent Realisation Of His Goal And Is Liberated, Thus Fulfilling The Ultimate Aim Of The Teaching. The Doctrine Of Vibration Is Not Meant For The Spiritually Dull. It Is Not For The Worldly Whose Consciousness, Clouded By Ignorance, Is As If Dreaming, Even During The Waking State Of Daily Life, The Dream Of Its Own Thought-Constructs. The Teachings Are Meant For Those Who Are Awake (Prabuddha), Those Who, Full Of Faith And Reverence, Are Always Alert And Intent On Discerning The True Nature Of Ultimate Reality. This Reality Is Understood In Three Basic Ways. The First Is Purely Transcendental. The Stanzas Choose This Aspect As The One

Which Formally Defines It Most Specifically. Ultimate Reality Transcends All The Opposites, Including Subject And Object. This Does Not Mean, However, That It Is An Unconscious Void, A Mere Absence Of All Existence. In Fact, This Negative Characterization Of Reality (Which Includes Also A Denial Of All That Is Unconscious) Implies A Positive Immanence In Which The Opposites Are United In The Oneness Of Pure Consciousness That Is Equally Siva And Spanda, His Universal Activity. These Two Seemingly Contrasting Aspects Are Reconciled In The Third, Namely, Reality Understood As The Essential Nature Of All Things. Although Universal And Everywhere The Same, It Is Understood To Be The Essential And Specific Nature Of Each Existent As Its 'Own Nature' (Svabhava).

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the

total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

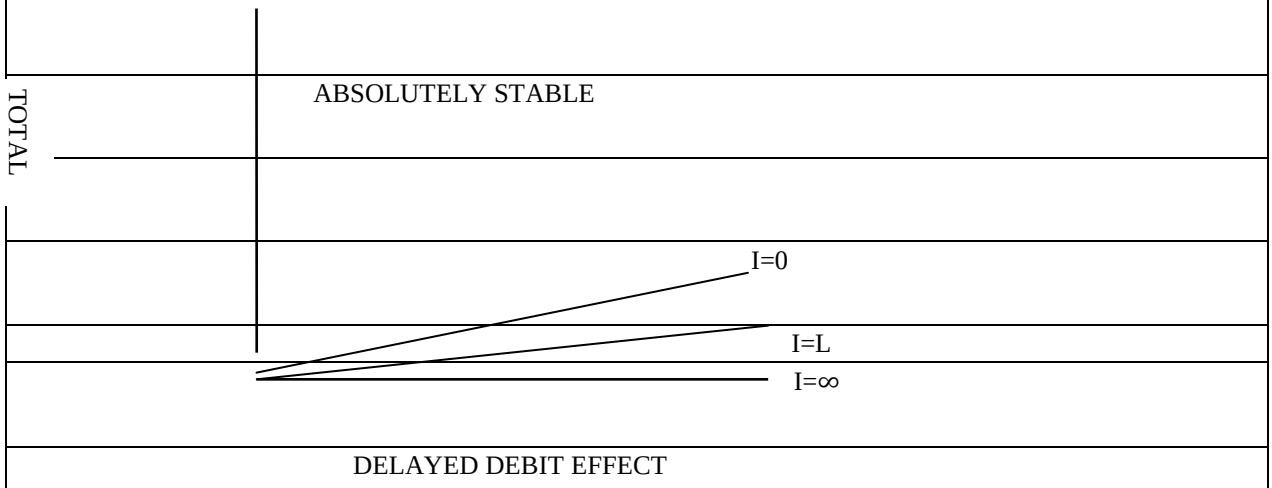



FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

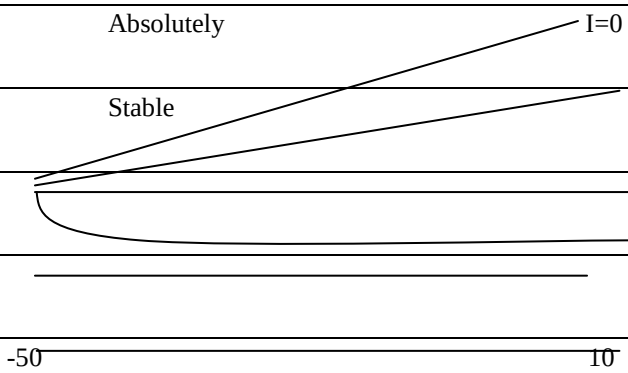
Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.

5

DEBIT REGULATION




**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

This equation is of paramount importance in the following cases: According to the principles of geometric quantum mechanics, the physical characteristics of a given quantum system can be represented by geometrical features that are preferentially identified in this complex manifold. Here we construct a number of examples of such features as they arise in the state spaces for spin, spin 1, spin and spin 2 systems, and for pairs of spin systems. A study is then undertaken on the geometry of entangled states. A locally invariant measure is assigned to the degree of entanglement of a given state for a general multi-particle system, and the properties of this measure are analysed for the entangled states of a pair of spin particles. With the specification of a quantum Hamiltonian, the resulting Schrödinger trajectories induce an isometry of the Fubini–Study manifold, and hence also an isometry of each of the energy surfaces generated by level values of the expectation of the Hamiltonian. For a generic quantum evolution, the corresponding Killing trajectory is quasiergodic on a toroidal subspace of the energy surface through the initial state. When a dynamical trajectory is lifted orthogonally to Hilbert space, it induces a geometric phase shift on the wave function. The uncertainty of an observable in a given state is the length of the gradient vector of the level surface of the expectation of the observable in that state; a fact that allows us to calculate higher orders corrections to the Heisenberg relations. A general mixed state is determined by a probability density function on the state space, for which the associated first moment is the density matrix. The advantage of a general state is in its applicability in various attempts to go beyond the standard quantum theory, some of which admit a natural phase-space characterization. **Journal of Geometry and Physics Volume 38, Issue 1, April 2001, Pages 19–53 Geometric**

**quantum mechanics Dorje C. Brody Lane P. Hughston DOI: 10.1016/S0393-0440(00)00052-8** Expansion (unmesa) of [consciousness] or the creative intuition {pratibha) is (=) [experienced] in the interval which divides two [moments] of differentiated perception (vikalpa) It is here that they arise and disappear. The Sastras and Agamas proclaim with reasoned argument that it is (=) free of thought-constructs {Nirvikalpa) and precedes (e) Path to Liberation all mental representations of any object. None can deny that a gap exists between perceptions insofar as two moments of thought are (=) invariably divided. This [gap] is (=) the undifferentiated unity of all the countless manifestations. Similarly, in the outer more objective sphere, where change consists (e) of the alterations in the configurations of manifest appearances (Abhasa\ the transition from one to (e&eb) another corresponds to (e&eb) a phase of pure luminosity that marks (eb) the beginning of one form and the end of another

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^\infty \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7 \left( U_7 - V_7 \int_0^\infty \cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7 \right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7 \int_0^\infty \cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$ then given the constraints of $f(T_7), (A1.5)$ *can never be satisfied* Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7\lambda_7 e^{-T_7\lambda_7} + [(U_7 - V_7)M_7 + Y_7] - [((U_7 -$$

$$V_7 M_7 -] e^{-T_7\lambda_7} = 0$$

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$Z_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7)\cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7)\sin(Z_7 T_7) = 0$$

The warrant of Truth is ever elusive when we deconstruct the foundational justifications of our convenient notions about the way the world works. It is easy to confuse what is actually the creation of beliefs with the "discovery of Truth," a common goal of science and theology. Science offers no ontological Scientific Picture of the World but substitutes a number of applied theories as meta-theory. It's close, powerful and

useful: it 'works' and allows use to do work. But the still mysterious "true" metatheory may reveal the basis of consciousness. But who knows when we will lift that veil? Permeating the living reality of our culture are certain contagious notions, fads and trends that have the ability to influence the way we think about the nature of Reality and ourselves. Some of them are toxic and can consume you; others are just whims of pop culture infectious and benign as the common cold. **Jung** described the concept of psychic contagion by certain archetypal forces inherent in the human psyche, which manifest in our spiritual lives and belief systems. He spoke of both conscious and unconscious contamination. Notions like this range from simple superstitions to scientific concepts, to urban myths. Notions sweep through our culture and insinuate themselves within its fabric, as fads, whether they are "real" or not, they can be influential. They mood alter us, make us feel we belong. An analysis of these notions is useful in distinguishing a common human phenomenon from any potential "alien influence" which may or may not be exerted on us from an unknown source. **Strange Attractors: Transference, Holography, and an Archetype Burke, J. (2003). Strange Attractors: Transference, Holography, and an Archetype (Doctoral dissertation, Pacifica Graduate Institute, 2003).** In consideration to the Nobel Proze for Chemistry this year(2014), a stability nalysis of the proble in to single molexues would be a contemporaneous question that has to be taken cognisace of and given importance to . Determination of molecular structure by geometry optimization became routine only after efficient methods for calculating the first derivatives of the energy with respect to all atomic coordinates became available. Evaluation of the related second derivatives allows the prediction of vibrational frequencies if harmonic motion is estimated. More importantly, it allows for the characterization of stationary points. The frequencies are related to the eigenvalues of the Hessian matrix, which contains second derivatives. If the eigenvalues are all positive, then the frequencies are all real and the stationary point is a local minimum. If one eigenvalue is negative (i.e., an imaginary frequency), then

the stationary point is a transition structure. If more than one eigenvalue is negative, then the stationary point is a more complex one, and is usually of little interest. When one of these is found, it is necessary to move the search away from it if the experimenter is looking solely for local minima and transition structures. The total energy is determined by approximate solutions of the time-dependent Schrödinger equation, usually with no relativistic terms included, and by making use of the Born–Oppenheimer approximation, which allows for the separation of electronic and nuclear motions, thereby simplifying the Schrödinger equation. This leads to the evaluation of the total energy as a sum of the electronic energy at fixed nuclei positions and the repulsion energy of the nuclei. Notable exceptions are certain approaches called direct quantum chemistry, which treat electrons and nuclei on a common footing. Density functional methods and semi-empirical methods are variants on the major theme. **Orlando Alvarez** discusses Polyakov's quantization of the string in the presence of a boundary allowing for an arbitrary topology for the world sheet. In addition to the dynamical conformal factor discovered by Polyakov, there are a finite number of new degrees of freedom if the surface is more complicated than a sphere or a disc. The quantization of the Liouville theory in an arbitrary topology is also discussed. A one-loop calculation shows that the model is renormalizable if one performs a mass renormalization and an additive field renormalization. The renormalization group equations have a perturbative infrared unstable fixed point in all topologies. Copyright © 1983 Published by Elsevier B.V. **Nuclear Physics B Volume 216, Issue 1, 25 April 1983, Pages 125–184 Theory of strings with boundaries: Fluctuations, topology and quantum geometry Orlando Alvarez DOI: 10.1016/0550-3213(83)90490-X** Quantum geometry (the modern loop quantum gravity involving graphs and spin-networks instead of the loops) provides microscopic degrees of freedom that account for black-hole entropy. However, the procedure for state counting used in the literature contains an error and the number of the relevant horizon states is underestimated. In our paper a

correct method of counting is presented by **Marcin Domagala and Jerzy Lewandowski.** Results lead to a revision of the literature of the subject. It turns out that the contribution of spins greater than 1/2 to the entropy is not negligible. Hence, the value of the Barbero–Immirzi parameter involved in the spectra of all the geometric and physical operators in this theory is different than previously derived. Also, the conjectured relation between quantum geometry and the black-hole quasi-normal modes should be understood again. **Marcin Domagala and Jerzy Lewandowski 2004 Class. Quantum Grav 21 5233 doi:10.1088/0264-9381/21/22/014 Black-hole entropy from quantum geometry** Quantum geometry predicts that a universe evolves through an inflationary phase at small volume before exiting gracefully into a standard Friedmann phase. This does not require the introduction of additional matter fields with ad hoc potentials; rather, it occurs because of a quantum gravity modification of the kinetic part of ordinary matter Hamiltonians. Authors draw the cognizance of application of the same mechanism that can explain why the present day cosmological acceleration is so tiny. DOI: http://dx.doi.org/10.1103/PhysRevLett.89.261301 **Inflation from Quantum Geometry Phys. Rev. Lett 89, 261301 – Published 12 December 2002 Martin Bojowald** The loop quantum cosmology of the closed isotropic model is studied with special emphasis on a comparison with traditional results obtained in the Wheeler-DeWitt approach by **Martin Bojowald** .Investigation includes the relation of the dynamical initial conditions to boundary conditions such as the no-boundary or the tunneling proposal and a discussion of inflation from quantum cosmology. **DOI: http://dx.doi.org/10.1103/PhysRevD.67.124023 Loop quantum cosmology, boundary proposals, and inflation Phys. Rev. D 67, 124023 – Published 23 June 2003 Martin Bojowald and Kevin Vandersloot** In loop quantum cosmology, the universe avoids a big bang singularity and undergoes an early and short super-inflation phase. During super-inflation, non-perturbative quantum corrections to the dynamics drive an inflaton field up its potential hill, thus setting the initial conditions for standard inflation.

**Shinji Tsujikawa et al** show that this effect can raise the inflaton high enough to achieve sufficient e-foldings in the standard inflation era. They also analyse the cosmological perturbations generated when slow-roll is **violated after super-inflation** and show that loop quantum effects can in principle leave an **indirect signature** on the largest scales in the CMB, with some loss of power and running of the spectral index. **Shinji Tsujikawa et al 2004 Class. Quantum Grav 21 5767 doi:10.1088/0264-9381/21/24/006 Loop quantum gravity effects on inflation and the CMB Quantum geometry from quantum information** is the study of **Florian Girelli and Etera R Livine .**Loop quantum gravity defines the quantum states of space geometry as spin networks and describes their evolution in time. We reformulate spin networks in terms of harmonic oscillators and show how the holographic degrees of freedom of the theory are described as matrix models. This allows us to make a link with non-commutative geometry and to look at the issue of the semi-classical limit of loop quantum gravity from a new perspective. This work is thought of as part of a bigger project of describing quantum geometry in quantum information terms. **Florian Girelli and Etera R Livine 2005 Class. Quantum Grav 22 3295 doi:10.1088/0264-9381/22/16/011 Reconstructing quantum geometry from quantum information: spin networks as harmonic oscillators.** Quantum geometric mechanics is the topic of study by **Dorje C. Brody Lane P. Hughston** establishing a clear relation between the quantum mechanics and quantum geometry and also the entangled states. The manifold of pure quantum states can be regarded as a complex projective space endowed with the unitary-invariant Fubini–Study metric. Note how differential gheometry concepts with Nunney's model could be entangle so that it is feasibly and viably applied to various type of model including consciousness (See Doctrine of Vibration)

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $\boldsymbol{Z}_7^2$ and $Z_7$in (A2.2):

$W_7^2(X_7^2-3M_7A_7X_7-M_7^2(X_7-Y_7)+2M_7^3A_7+2A_7^2M_7^2]$

$$+W\left[MA^3-A^2\mathrm{X}+2\mathrm{MA}(2\mathrm{X}+\mathrm{Y})-2X^2\right]+A^2(\mathrm{X}+\mathrm{Y})=0$$

*where* $W_7$=1-cos ($\mathrm{Z}_7T_7$).

In the absence of a delay the system is stable for $A_7>0,$ and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7>(\mathrm{X}_7-Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$A_7>[\ \mathrm{X}_7-Y_7/2\ M_7$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_7,A_7)=0$$

Which is quadratic in $W$ the shape of the quadratic function in the range 0⍰$W_7$ ⍰2 *is* of interest First consider the extreme values

$Q_7(0,A_7)=A_7^2\ (\mathrm{X}_7-Y_7)$

$Q_7(2,A_7)=(A_7+2\ \mathrm{M}_7)^2(2A_7\mathrm{M}_7-(\mathrm{X}_7-Y_7))$

$which$ are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7M_7^2 - 3M_7X_7\right) + W_7\left(3M_7A_7^2 - 2A_7X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7 + Y_7\right)$$

$$=2M_7^2W_7^2(M_7 + A_7) + M_7A_7^2W_7 + \left(2A_7M_7 - X_7 + Y_7\right)(A_7W_7 + M_7W_7^2$$

$$+ (2- W_7)((2M_7X_7 + M_7Y_7)W_7 + A_7(X_7 + Y_7)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $W_7$touches the $W$ –axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$has no minimum in the range 0≤ $W_7$ ≤2.Defining

$Q_7\left(W_7, A_7\right) = a_7W_7^2 + c_7$ $where\ a_7\ b_7\ and\ c_7$ Are functions of $A_7$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7 ≥ b_7 ≥0$$

And

$$b_7^2 ≤4a_7c_7$$

$$Case(i): Y_7 = 0$$

$Let\ Y_7 = 0$ And try the substitution A= $X_7 = /2M_7\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_7 X_7 / 2M_7$) is positive for *W<2. Since* $Q_7$ increases with $A_7$(see A2.9), increasing $A_7$will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

$Case\ (ii): Y_7 = y_7\ X_7$ ,

$Let\ Y_7 = \ y_7\ X_7$And try the substitution $A_7 = X_7/(2 + A_7\ y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + \ n_7 y_7)$ we have

$$W_7^2\ y_7\left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2\left(2 - n_7 + n_7 y_7\right)\right]$$

$$-W_7\ X_7\left[\frac{X_7(1 + n_7 y_7)}{M_7^2(2 + n_7 y_7)^2} + 2\ y_7\left(n_7 - 1\right) + \frac{X_7^2(7 + y_7)}{M_7^2(2 + n_7 y_7)}\right] = 0.\ (A2.12)$$

Since $n_7$ and $A_7$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7), the\ condition(A2.77?)can\ be\ written\ as$

$\left[\frac{X_7(E_7 - 7)}{E_7^2}\right]^2\left[\frac{X_7^2}{M_7^2} - 4\ y_7\ E_7^2\frac{X_7}{M_7^2} + 4y_7(n_7 - 2)E_7^3\right]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$$2E_7[y_7E_7 + (y_7E_7(7+y_7)(E_7-n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7[y_7E_7 - y_7E_7(7+y_7)(E_7-n_7))^{7/2}]$$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2$) is large or small then (A2.6) will have a real solution although if may not be in the range (0 W 7 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7-7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2 - 2y_7\frac{(E_7+n_7^2y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right) - 4y_7(E_7-n_7)$$

And has the solution

$$\frac{X_7}{M_7^2} > \frac{E_7}{(E_7-7)}[y_7(E_7+n_7^2y_7) + [y_7^2(E_7+n_7^2y_7)^2+4y_7(E_7-7)(E_7-n_7)]^{7/2}]$$

Given $E_7 > n_7$this lower limit to $X_7 / M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7 / M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$in the

appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7 < n_7$ $(ie., 2A_7 M_7^2 < X_7 - Y_7)$then (A2.7)and (A2.8)show that $Q_7$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - (y_7 E_7(7 + y_7)(E_7 - n_7)]^{7/2}]$$

And

$$E_7 > n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2 > 2Y_7$then $A_7 > (X_7 - Y_7)/2M_7$.

If $A_7^2 < 2Y_7$then $A_7 > (X_7 - Y_7)/2M_7$.+ $(Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$and $X_7$and $Y_7$by $W_7^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7 \lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]e^{2T_7 \lambda_7} + M_7 V_7 e^{2T_7 \lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i}\, Z_7$ *and* separate (A3.1)into real and imaginary parts

$$Z_7^2 + \left(M_7 + V_7\right) Z_7 \sin\left(Z_7 T_7\right) - \left(U_7 M_7 + Y_7\right) + \left(M_7\left(U_7 + V_7\right) - X_7\right)$$

$$\mathrm{X}\cos\left(Z_7 T_7\right) - 2M_7 V_7 \cos^2\left(Z_7 T_7\right) + M_7 V_7 = 0$$

$$\left(M_7 + U_7\right) Z_7 - \left(M_7 + V_7\right) Z_7 \cos\left(Z_7 T_7\right) + \left(M_7\left(U_7 + V_7\right) - X\right)$$

$$\mathrm{X}\sin\left(\left(Z_7 T_7\right) - M_7 V_7 \sin\left(Z_7 T_7\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_7$and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - \cos Z_7 T_7$and $A_7 = U_7 - V_7$we obtain

$$4 M_7^2 V_7^2 W_7^4 + 2 M_7 V_7 W_7^3 [(M_7 - V_7)^2 + A_7\left(M_7 + V_7\right) + 2(A_7 M_7 - X_7)]$$

$$+ W_7^2 \left[M_7 - V_7\right)^2 \left(2 A_7 M_7 + Y_7 - X_7\right) + \left(2A_7 M_7 - X_7\right)^2 - A_7\left(M_7 - V_7\right)\left(2A_7 M_7 - X_7\right)$$

$$+ M_7 V_7 (A_7^2 + 4\left(X_7 + Y_7\right))]$$

$$+ W_7 \left[A_7^2 + 2X_7\right)\left(2 A_7 M_7 - X_7\right) + 2A_7 Y_7 (M_7 + V_7) - M_7 A_7^3] + A_7^2\left(X_7 M_7 - Y_7\right) = 0$$

Being a quartic in $W_7$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7 > 0$$

$$and$$

$$C_7 > (X_7 - Y_7)/2M_7$$

$recalling\ that\ A_7 = U_7$ and defining $C_7 = U_7 + V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0, U_7 V_7 = A_7^2 (X_7 + Y_7)$$

$$Q_7\left(2, U_7 V_7\right) = \left(C_7 + 2M_7^2\right)^2 (2C_7 M_7 + Y_7 - X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$ the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

## Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

**Module numbered five**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in

the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \quad (2)$$

This is the condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

## DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\left[F\left(R(t)\right) - M\left(R(t)\right)\right]$$

(3)

$$\frac{dR(t)}{dt} = B\left(R(t-T)\right) - D\left(R(t)\right) - N(t)G\left(R(t)\right).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t-T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\left(R(t)\right)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the

TIME DELAY

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

As it is usually expressed outside of phenomenological texts, to undergo a conscious experience necessarily means that there is something it is like for the subject to have that experience (Nagel 1974; Searle 1992). This is obviously true of bodily sensations like pain. But it is also the case for perceptual experiences, experiences of desiring, feeling, and thinking. There is something it is like to taste chocolate, and this is different from what it is like to remember what it is like to taste chocolate, or to smell vanilla, to run, to stand still, to feel envious, nervous, depressed or happy, or to entertain an abstract belief. Yet, at the same time, as I live through these differences, there is something experiential that is, in some sense, the same, namely, their distinct first-personal character. All the experiences are characterized by a quality of mineness or for-me-ness, the fact that it is I who am having these experiences. All the experiences are given (at least tacitly) as my experiences, as experiences I am undergoing or living through. All of this suggests that first-person experience presents me with an immediate and non-observational access to myself, and that consequently (phenomenal) consciousness consequently entails a (minimal) form of self-consciousness. To put it differently, unless a mental process is pre-reflectively self-conscious there will be nothing it is like to undergo the process, and it therefore cannot be a phenomenally conscious process. The mineness in question is not a quality like being scarlet, sour or soft. It doesn't refer to a specific experiential content, to a specific what; nor does it refer to the diachronic or synchronic sum of such content, or to some other relation that might obtain between the contents in question. Rather, it refers to the distinct givenness or the how it feels of experience. It refers to the first-personal presence or character of experience. It refers to the fact that the experiences I am living through are given differently (but not necessarily better) to me than to anybody else. It could consequently be claimed that anybody who denies the for-me-ness of experience simply fails to recognize an essential constitutive aspect of experience. Such a denial would be tantamount to a denial of the first-person perspective. It would entail the view that my own mind is either not given to me at all — I would be mind- or self-blind — or is presented to me in exactly the same way as the minds of others. There are also lines of argumentation in contemporary analytical philosophy of mind that is close to

and consistent with the phenomenological conception of pre-reflective self-awareness. Alvin Goldman provides an example: [Consider] the case of thinking about x or attending to x. In the process of thinking about x there is already an implicit awareness that one is thinking about x. There is no need for reflection here, for taking a step back from thinking about x in order to examine it…When we are thinking about x, the mind is focused on x, not on our thinking of x. Nevertheless, the process of thinking about x carries with it a non-reflective self-awareness (Goldman 1970, 96). A similar view has been defended by Owen Flanagan, who not only argues that consciousness involves self-consciousness in the weak sense that there is something it is like for the subject to have the experience, but also speaks of the low-level self-consciousness involved in experiencing my experiences as mine (Flanagan 1992, 194). As Flanagan quite correctly points out, this primary type of self-consciousness should not be confused with the much stronger notion of self-consciousness that is in play when we are thinking about our own narrative self. The latter form of reflective self-consciousness presupposes both conceptual knowledge and narrative competence. It requires maturation and socialization, and the ability to access and issue reports about the states, traits, dispositions that make one the person one is. To claim that every kind of self-consciousness is conceptual is overly cognitive. Bermúdez (1998), to mention one further philosopher in the analytic tradition, argues that there are a variety of **nonconceptual forms of self-consciousness** that are "logically and ontogenetically more primitive than the higher forms of self-consciousness that are usually the focus of philosophical debate" (1998, 274; also see Poellner 2003). This growing consensus across philosophical studies supports the phenomenological view of pre-reflective self-consciousness. (References: Stanford Encyclopedia, Post Modersnist Theories, Internet Encyclopedia of Philosophy) Korzybski pointed out that there are two ways to slice easily through life; to believe everything or to doubt everything. Both ways save us from thinking. So, from time to time you have to give yourself a Virtual Reality Check. The hallmark of postmodern philosophy has been disbelief or skepticism of all "metanarratives," or translations of reality. Postmodernism has even turned its profound skepticism on such important humanist concepts

as "objective truth" and reason.Yet, for a deconstructionist postmodern society, individually we are still riddled with superstition and gullibility, and open to manipulation through our belief systems as any politician, philosopher, clergy, or salesperson will attest. Further, most people are painfully naive when it comes to even the simplest scientific understanding. Most of us don't have a clue about the fundamental nature of physical reality or our own psychological nature, and our ability to be fooled by our senses and mind, or methods of social and technological persuasion. We may think we know, but the penetration is deeper than we dare imagine. Nature rejects the naiveté that seeks absolute truth. We are beginning to realize, individually and culturally, that "realities" are all human constructions. The task becomes one of "catching ourselves in the act" of creating our own "reality" from the flow of events. Human truth is always an engagement of mind with experience.. We don't need to fear the collapse of our personalistic belief system (the "box" we live in), nor our belief in absolute truth. **A strong desire to engage in the "quest for uncertainty" complements our anxiety that perhaps there is no absolute, objective ground to reality**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t - T_8)F_8(R_8(t - T_8) - N_8(t)M_8((R_8(t) \qquad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8}+\frac{N_8 dG_8}{dR_8}-\frac{dB_8}{dR_8}>\frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8}+\frac{dM_8}{dR_8}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_8=\frac{N_8 G_8}{2M_8}\frac{d(8nF_8)}{dR_8}(\geq 0) \quad (8a)$$

$$I_8=-\frac{N_8 G_8}{2M_8}\frac{d(8nM_8}{dR_8}(\geq 0) \quad (8b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8>L_8-I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_8>L_8-I_8$, when $P_8^2/4>I$ (8a)

$P_8>L_8-I_8+\frac{(I_8-P_8^2}{(L_8-I)}$ when $P_8^2/4<I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

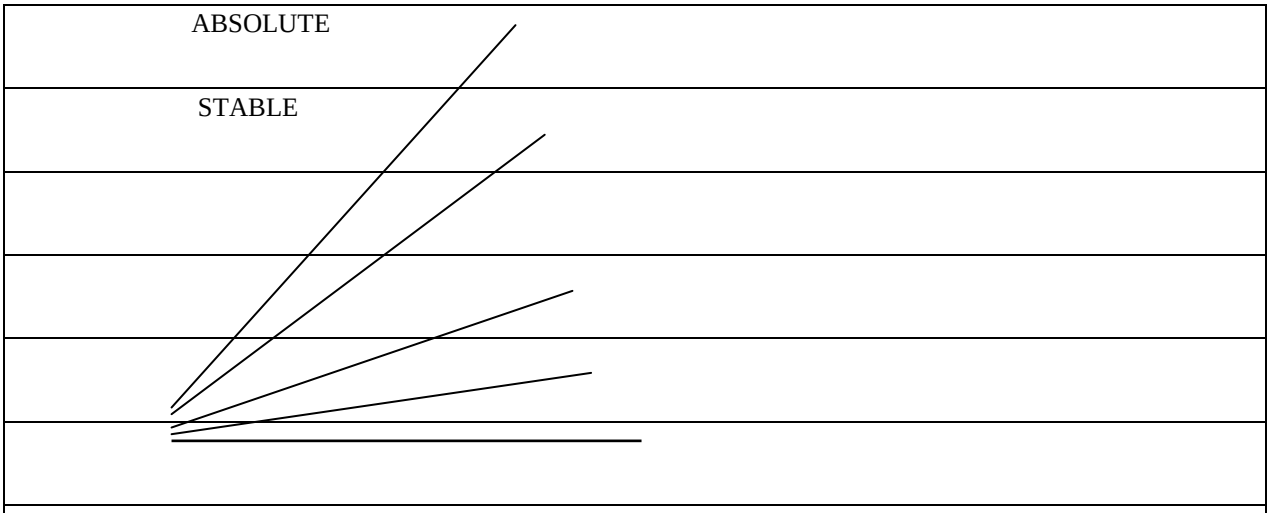



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

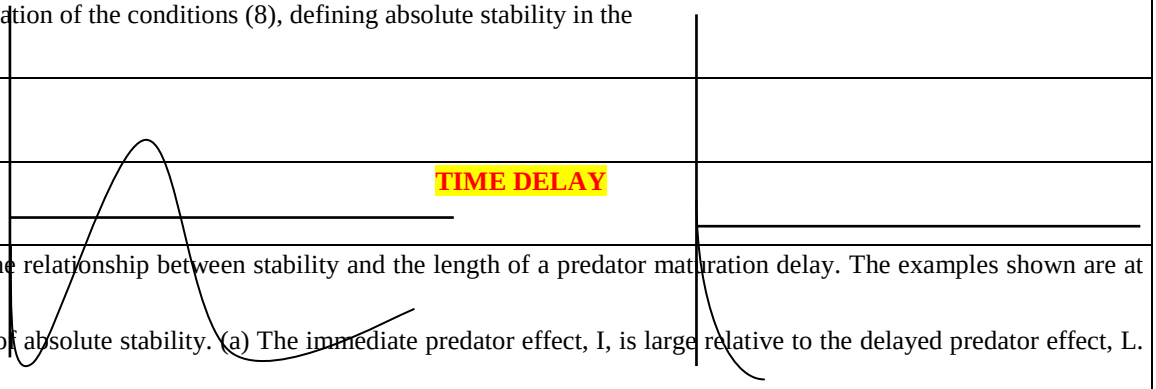



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior , we must note that the model could be applied to psychological systems too, albeit we are still not sure of the fact how the wave function collapses thereby due to perception.

Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility. He privatizes the sublime, having learned from the fate of his predecessors that the public can never be more than beautiful."(11) Notice is made here of what is indeed the telling feature. Levinas is praised, for example, for his alarm when "the social will [is] sought in an ideal of fusion . . . the subject losing himself in a collective

representation, in a common ideal. . . . It is the collectivity which says 'us,' and which, turned toward the intelligible sun, toward the truth, experience, the other at his side and not face to face with him. The tellingly titled D'un ton apocalyptique contains passages that are particularly excited: These people situate themselves outside the ordinary, but they have in common this: they describe themselves as having an immediate and intuitive relation with mystery. And they want to attract, to seduce, and lead others to the mystery, through mystery. This agogic function of the leader of men, of the duce, of the Führer, of the leader, places him above the crowd that he manipulates with the aid of a small numbers of adepts joined together in a sect with a secret language, a clique or a small party with its ritualized practices. The mystifiers pretend to have exclusive access to the privilege of a secret mystery. . . . The revelation or the unveiling of the secret is something that they jealously reserve for themselves. Jealousy is here a major characteristic. There is not narcissism and non-narcissism; there are narcissisms that are more or less comprehensive, generous, open, extended. What is called non-narcissism is in general but the economy of a much more welcoming, hospitable narcissism, one that is much more open to the experience of the other as other. I believe that without a movement of narcissistic reappropriation, the relation to the other would be absolutely destroyed; it would be destroyed in advance. The relation to the other--even if it remains asymmetrical, open, without reapporpriation--must trace a movement of reappropriation in the image of oneself for love to be possible, for example. Love is narcissistic. Beyond that, there are little narcissisms, there are big narcissisms. . . .(23) **Justice of the Pieces Deconstruction as a Social Psychology Douglas Collins** Sabine Wilke cogently argues for the rise of the aesthetic dimension in philosophical writing by showing how the metaphorical and poetic style of Heidegger, the principles of phrase configurations with Adorno, or the topographic writing of derrida replace through their very form something what traditionally would have been subject to discursive exposure The concern shifts from telling to showing, showing first and foremost that something remains untold. An exemplary book in this aesthetic drive is Avital Ronell's Telephone Book where the playful layout and explicit break with traditional design interrupt the

sequential logic of the paragraphs thereby performing the very argument at stake: open up new circuits of meaning and signification. The engagement with form and publishing formats is only consequential since media theory as exposed by Walter J. Ong, Eric A. Havelock, Parry Miman or Marshall McLuhan besides many others made it clear that different communication technologies condition the mode of thought of a given culture The fractals of Benoit Mandelbrot or the butterfly shaped Lorenz attractor are the popular icons of this turn in the sciences. For communication studies and philosophy, Vilém Flusser was one of the pioneer thinkers reflecting the rise of what he called "techno-images." The photograph is the paradigmatic case which condenses the whole crises of Western thought and culture on its grainy surface. It subverts the linear logic of writing.6 Flusser's central diagnosis detects a shift from the grand narratives associated with the literate mind and Modernity towards creating and envisioning the world through image based practices. Writing himself in a succinct, visually inspired style, Flusser contents that the web of meaning, spun throughout the centuries by our textual practices, has become transparent. We discover the void between the letters through which the world escapes: "We are alienated from the world of text, precisely because we see through them as our own product."7 As a consequence we are confronted with a world of dissociated facts and information bits. Flusser hails this fragmented world as the advent of the zero dimensional point-universe that will provide the new raw material for our image based practices. Techno-images are nothing else than a conscious, inherently artistic drawing-together of point-elements as exemplified first and foremost by the silver-grains of the photograph but also by more recent pixel based techno-images. Similar to blueprints, designs or statistical curves, the weather map for example visualizes and contracts a tremendous amount of meteorological data from widely distributed stations into a single image of high- and low pressure

areas, temperatures, wind speeds, etc. Techno-images "compute" in this sense the universe of unimaginable, abstract information-bits into a concrete surface – the image. Echoing Paul Feyerabend, these image practices have to be understood as creative practices where the truth of scientific models has been replaced by its aesthetic quality

**Aesthetics, Writing, Networked Computers Jörg Muller Dissertation Submitted to the Division of Media and Communications of The European Graduate School in Candidacy for the Degree of Doctor of Philosophy**

**Lothar Schäfer (1) and Sisir Roy (2) in** Quantum Reality, the Importance of Consciousness in the Universe, the Discovery of a Non-Empirical Realm of Physical Reality, and the Convergence with Ancient Traditions of Indian and Western Philosophy investigate metaphysical aspects that include 1) the discovery of a non-empirical part of physical reality in a realm of potentiality; 2) the emanation of the empirical world out of a realm of non-material forms; 3) the discovery that the nature of physical reality is that of an indivisible Wholeness – the One; and 4) panpsychism: the possibility that the One is aware of its processes like a Cosmic Consciousness. The convergence of powerful traditions of seemingly disparate cultures is particularly important to point out in the present process of globalization, when a unifying view is needed to avoid controversy and conflict. To quote them in extenso : This is the nonlocality of the quantum world. Menas Kafatos and Robert Nadeau (1990) have drawn a remarkable conclusion from this phenomenon: If reality is nonlocal, the nature of the universe is that of an undivided wholeness. Because our consciousness has emerged from this wholeness and is part of it, it is possible to conclude that an element of consciousness is active in the universe: a Cosmic Consciousness. Pre-reflective self-consciousness is pre-reflective in the sense that (1) it is an awareness we have before we do any reflecting on our experience; (2) it is an implicit and first-order awareness rather than an explicit or higher-order form of self-consciousness. Indeed, an explicit reflective self-consciousness is possible only because there is a pre-reflective self-awareness that is an on-going and more primary self-consciousness. Although phenomenologists do not always agree on important questions about method, focus, or even whether there is an ego or self, they are in close to unanimous agreement about the idea that the experiential dimension always involves such an implicit pre-reflective self-awareness. In line with Edmund Husserl (1959, 189, 412), who maintains that consciousness always involves a self-appearance (Für-sich-selbst-erscheinens), and in agreement with Michel Henry (1963, 1965), who notes that experience is always

self-manifesting, and with Maurice Merleau-Ponty who states that consciousness is always given to itself and that the word 'consciousness' has no meaning independently of this self-givenness (Merleau-Ponty 1945, 488), Jean-Paul Sartre writes that pre-reflective self-consciousness is not simply a quality added to the experience, an accessory; rather, it constitutes the very mode of being of the experience: This self-consciousness we ought to consider not as a new consciousness, but as the only mode of existence which is possible for a consciousness of something (Sartre 1943, 20 [1956, liv]). The notion of pre-reflective self-awareness is related to the idea that experiences have a subjective 'feel' to them, a certain (phenomenal) quality of 'what it is like' or what it 'feels' like to have them.

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the

total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

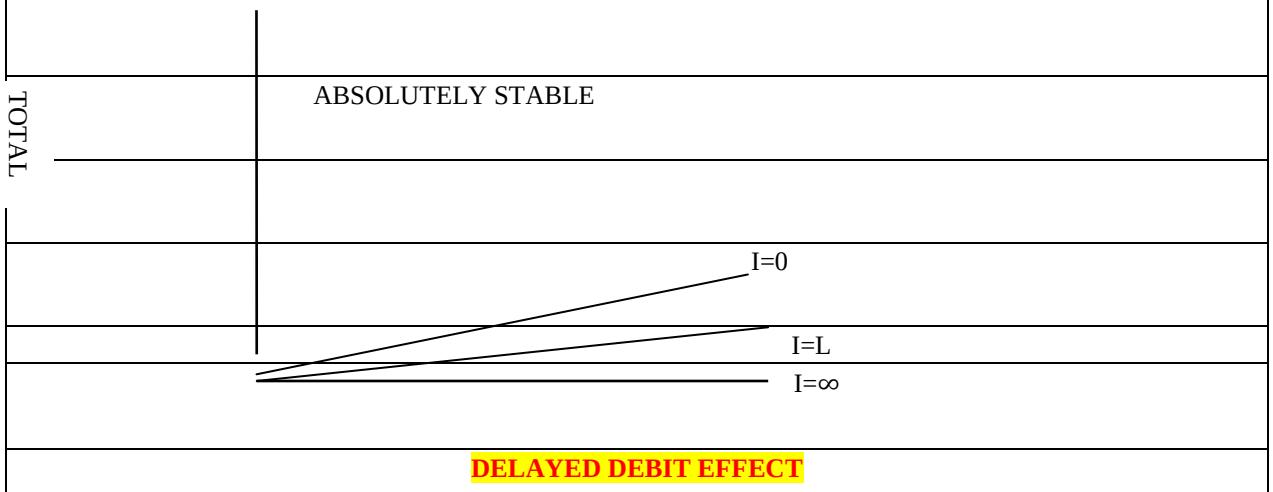
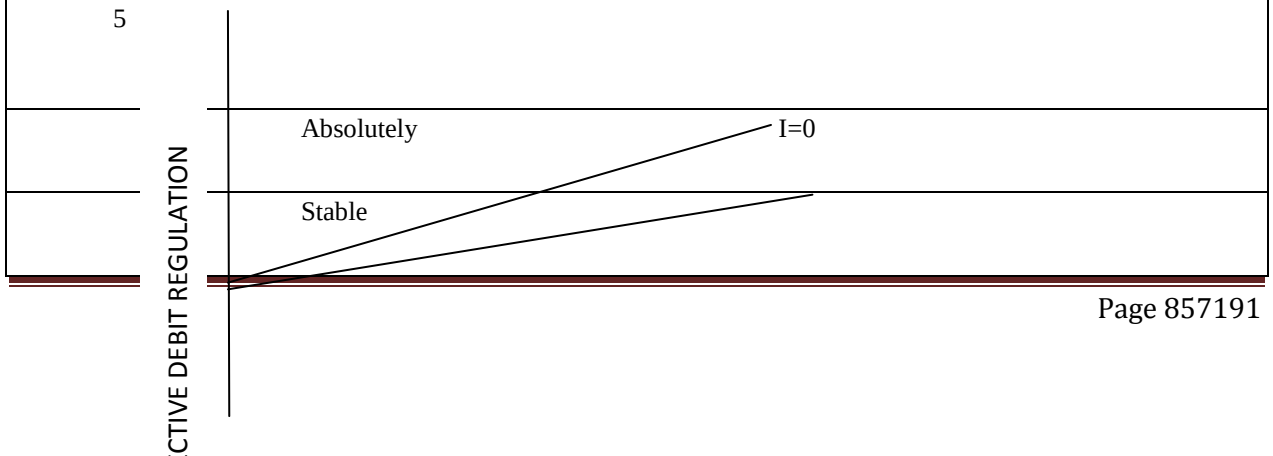


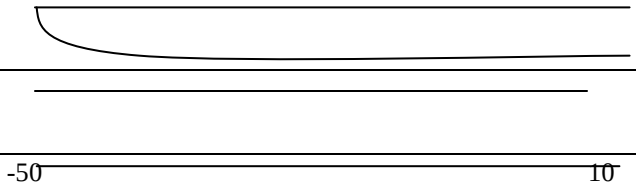

FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.






DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

**Utility Analysis And neural Network**

Caroline Williams investigated affective density of the political and its effect on our understanding of political subjectivity. Taking up Spinoza's challenge to think about affect beyond corporeal embodiment, he argues that there is a modality of affectivity that cannot simply be inscribed within the borders of subjectivity. theorising affect as an impersonal force anchored in a relational ontology that gives due recognition to the circulation of affects, as well as to their ambivalent structure in creating sites of identification, and he utilises this ontology to reflect on the dynamic of the political and the shape of political subjectivity. Williams argues that Spinoza's philosophy (through ideas of conatus and imagination) offers the conceptual resources to reconfigure the composition of affective subjectivity as a transindividual social bond and as an unconscious dynamic of ethico-political existence. **Subjectivity (2010) 3, 245–262. doi:10.1057/sub.2010.15 Affective processes without a subject: Rethinking the relation between subjectivity and affect with Spinoza Caroline Williams** Phenomenological Time bears ample evidence and infallible observatory and apodictic evidence albeit with a sense of chagrined circumvention and defeated disgruntlment circularity of both psychological and phenomenological interpretations of art which are sometimes insightful, but most often work on the logic of cyclical time and unavoidable drives and representations played out as a ‘natural’ order: where ‘men can’t avoid being boys’ a circularity which valorizes a helplessness/truthfulness conundrum. One begins with a classical phenomenological fallacy with Duchamp’s Fountain, does the urinal keep returning or do we keep returning to the urinal? Put in this way, it is a male question, something that

phenomenology conveniently conceals when positing the "I", or indeed, the "we". But we must be careful not to simply assume that the phenomenological first person is white, heterosexual, able-bodied and male, providing a foundation upon which knowledge should be based as a universal condition. Deleuze's second critique of time is linked to Kant. This model of time extends not only to phenomenology but to the phenomenological interpretation of art and is seen as a straight line. In the Critique of Pure Reason, Kant cuts up circular time and reassembles it as a series of sensory experiences which memory processes and which constitute subjectivity. Whereas the circular model of time (and interpretation) inscribes fatalism into the subject and the world, the phenomenological method is in danger of solipsism, where the subject constitutes the world. Much of embodied philosophy and enactivist theories of cognition go down this line with a plethora of art historical analyses by the likes of Rosalind Krauss, Michael Fried and others taking the lead from Merleau-Ponty. These approaches successfully challenge the **homunculus model of a Cartesian mind-body** dualism where there is a supervisor in the mind who processes experience and even consciousness and to whom all representations are directed. But phenomenology is mainly concerned with perceptual consciousness, the processing of colors, shapes, forms, motion in art and the world. The senses and sensorimotor processes are the fundamental pathways by which the interior is connected to the exterior in a chiasm of self-constitution in the world. The problem with this is that we explore the art through our sense perceptions, and what does art reveal to us?—our sense perceptions! Such solipsistic approaches have real problems relating complex and abstract conceptual production to bodily and sensorimotor contingencies, and they nearly always posit the subject as the site and object of discovery in the world. To be fair, phenomenology has moved on to rather more sophisticated models of intersubjectivity, taking into account the specificities of different kinds of bodies and ways of being beyond the white, male heterosexual as the common denominator, yet many naïve approaches to art perpetuate the notion that perceptual, rather than conceptual processes and evolutionary neural patterns are primary, fundamental ways to understand the art experience In fact, these aspects are only part of art

production and reception, aspects which are often filtered out or made known to us in order for us to parody such trigger responses, as in much of Duchamp's work, which holds in contempt such retinal approaches to art. If we take the **subject of abstract expressionism** in the 1950s or later minimalist art, art history, teaching and interpretation still insists on putting the act of perceiving as fundamental to the constitution not only of the phenomenological first person but also of art itself. The male act of usually standing up, which consists of directing a stream with the hand(s) into a container or on a flat surface, sand or snow in order to leave patterns and traces (the usual child's prank) underlies the logic of macho 1950s Abstract Expressionist culture, especially that of Jackson Pollock, so effectively parodied by Carolee Schneeman's Vagina Scroll, 1975. Here, painting or writing is a kind of a discharge which invests the **object and space with** the phenomenological "I", disciplined by existential struggle and reflected back by the traces, signatures and gestures left behind in the work of art, rediscoverable as an element in the making of other subjectivities.It is precisely this kind of solipsism which is blind to the art event as something well in excess of re-establishing habits of subjective self-identity through the production and interpretation of art. The attempt to reinforce a phenomenological identity upon the event by personalizing sense experience is unable to see art as something which remodels the phenomenological 'I' and neural plasticity itself. We learn new technologies and redraw our neural pathways through such new and spontaneous experiences which draw us into intersubjective worlds. Duchamp's Fountain and Deleuze's Repetition and Différence Grégory Minis sale (I have made some changes) The relation of different groups to deconstruction is settled by its relation to groups. An American enthusiast--one who has had a role in the development of our affection for Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^{\infty} \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$then given the constraints of $f(T_8)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8\lambda_8 e^{-T_8\lambda_8} + [(U_8 - V_8)M_8 + Y_8] - [((U_8 -$$

$$V_8M_8-]e-T_8\lambda_8=0$$

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$Z_8^2 + M_8Z_8\sin(Z_8T_8) - (A_8M_8 + Y_8) + (A_8M_8 + X_8)\cos(Z_8T_8) = 0$$

$$(A_8M_8)Z_8 - M_8Z_8\cos(Z_8T_8) + (A_8M_8 - X_8)\sin(Z_8T_8) = 0$$

Where $A_8 = U_8 - V_8$Using (A2.3), substitute for $Z_8^2$ and $Z_8$in (A2.2):

$W_8^2(X_8^2$-3 $M_8\ A_8\ X_8$- $M_8^2(X_8$- $Y_8) + 2\ M_8^3\ A_8$+2 $A_8^2\ M_8^2]$

$$+W\left[MA^3 - A^2X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_8$=1-cos $(Z_8T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8/2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \leq W_8 \leq 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2 (X_8 - Y_8)$$

$$Q_8(2, A_8) = (A_8 + 2M_8)^2 (2A_8 M_8 - (X_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2\left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8\left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8\left(2X_8 + Y_8\right)\right)$$

$$+2A_8\left(X_8 + Y_8\right)$$

$$=2M_8^2 W_8^2 (M_8 + A_8) + M_8 A_8^2 W_8 + (2A_8 M_8 - X_8 + Y_8)(A_8 W_8 + M_8 W_8^2$$

$$+ (2- W_8)((2M_8 X_8 + M_8 Y_8) W_8 + A_8 (X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $W_8$touches the $W$–axis only at its minimum value in which case (A2.9) shows that

further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range 0≤ $W_8$ ≤2.Defining

$Q_8(W_8, A_8) = a_8 W_8^2 + c_8$ *where* $a_8\, b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \geq b_8 \geq 0$$

And

$$b_8^2 \geq 4a_8 c_8$$

$$Case(i): Y_8 = 0$$

*Let* $Y_8 = 0$ And try the substitution A= $X_8 = /2M_8\, in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8$ / $2M_8$) is positive for *W<2. Since* $Q_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\, X_8$ ,

$Let\ Y_8 = y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 + n_8 y_8)$ we have

$$W_8^2\ y_8\left[X_8 n_8 \frac{(1+n_8 y_8)}{(2+n_8 y_8)} + M_8^2\left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8\ X_8\left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2} + 2\ y_8\left(n_8 - 1\right) + \frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)} = 0.\ (A2.12)\right.$$

Since $n_8$ and $A_8$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8 = (2 + n_8 y_8), the\ condition (A2.88?) can\ be\ written\ as$

$[\frac{X_8(E_8-8)}{E_8^2}]^2[\frac{X_8^2}{M_8^2} - 4\ y_8\ E_8^2\frac{X_8}{M_8^2} + 4y_8(n_8 - 2)E_8^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8 / W_8^2)$ can be expressed as

$2\ E_8[y_8 E_8 + (y_8\ E_8(8 + y_8)(E_8 - n_8))^{8/2}] > \frac{X_8}{M_8^2} > 2E_8[y_8 E_8 - y_8 E_8(8 + y_8)(E_8 - n_8))^{8/2}]$

Thoughts on the Social and Political Implications of Correlationism .Lion Mirror Truth be told, as my thought has evolved the issue of correlationism had fallen off the radar for me. Somehow the debate had come to seem too “philosophical” to me, too “scholastic”, too remote from what interests me: understanding why social assemblages are organized as they are, how power functions in social assemblages, and what we might do to address that power and change things. Somehow the question of whether or not we can get out of the correlation between thinking and being just came to seem remote from these sorts of issues. Somehow it

seemed too epistemological. ASIDE: Numerous discussions over the years have led me to believe that the debate over correlationism is poorly understood (or maybe I just don't understand it). On countless occasions I've heard people say "of course we must relate to things in order to know them." Well yeah, of course! I don't think this is what the critic of correlationism is getting at. It seems to me that correlationism is something more robust than the theses that we must relate to something to know it. Correlationism instead seems to require the theses that thought and being are indiscernible. Put more concretely, the correlationist is someone who argues that we either a) can never tell whether being is merely a construction of our thought (weak correlationism), or b) who argues that thought actually constructs being (strong correlationism). In other words, correlationism is another name for idealism. One can hold that we must relate to something in order to know it without being a correlationist. As an aside I should also add that I am a correlationist about some things. For example, I think money is something constructed by society and is therefore a strong correlationist when it comes to money. At any rate, for a long time I'd become rather indifferent to debates about correlationism and philosophies of access. I had learned the lessons of speculative realism– which I could have also learned, I think, from Deleuze and Guattari, the new materialist feminists, actor-network theorists such as Latour –and had moved on. However, occasionally you come across a tone of phrase that pitches something in a different light. In The Cut of the Real, Katerina Kolozova writes, …the political problem of contemporary philosophy identified by the 'new realists' is, in fact, the product of a more fundamental epistemic problem. In his book After Finitude, Quentin Meillassoux calls this problem 'correlationism' and identifies it as an essentially post-Kantian legacy, which continues to dominate and limit philosophy. As a matter of fact, correlationism lies at (eb) the heart of postmodern theory and consists in the premise that thought can only 'think itself,' that the real is inaccessible to knowledge and human subjectivity, and that there is nothing but discursive

constructs that fully determine thinking and that are meth0dologically accounted for all the way down. (1 – 2) Thought thinking only itself: Thought only encountering itself. In the jargon of postmodern and poststructuralist lingo, this would be the thesis of infinite semiosis, where signs ("thoughts") only ever relate to other signs. Within this framework, discursivity comes to be the hegemonic framework defining all of being. At the level of politics and social theory more generally, if the correlationist thesis is true the consequences are clear: all social phenomena are discursive and all solutions to social and political problems will be discursive. The sole sphere of the political will be the discursive and all questions of politics will be questions of speech-acts and interpretation. The problem here is not that many theorists recognize that the discursive and semiotic plays an important role in the social and the political. It does and I've repeated this tirelessly. The problem is with what happens when thought or the semiotic becomes a hegemon, an "all", foreclosing our ability to recognize other forms of power. What I've wanted to say is that not all power functions discursively. In my last post and elsewhere I spoke of some other forms of politics: Thermopolitics: The politics surrounding **energy in the form of calories and fuels such as gasoline** and coal, and how our life and our very bodies are structured by energy dependencies and by being trapped in particular distributive networks that render these forms of energy available. I'm being quite literal when I speak of energy, talking about the effects, for example, of the absence of food in certain educational environments on cognition, for example; and am generally hostile to metaphorical extensions of the concept of energy which I see as erasing the dimension of real materiality. Geopolitics: The role that features of natural and built geography such as mountain ranges, rivers, oceans, soil conditions, roads, housing design, etc., play in the form that social relations take and how they impact individual bodies. Chronopolitics: The way in which the structurations of time organize what is possible for us. For example, the structuration of the working day, how much we can say and comprehend at any given time, the impact

of things like the invention of the clock, etc. **Oikopolitics: This would be the domain of political economy described so well by Marxists**. So five different types of politics: Semiopolitics (or what currently dominates critical theory), thermopolitics, geopolitics, chronopolitics, and oikopolitics. No doubt there are other sites of the political or political struggle that we could speak of, but this is a good start. Also, it should be obvious that these aren't exclusive domains, but are entangled in all sorts of important ways. For example, something might take place at the level of semiopolitics (speech, law, rhetoric, norms, and communication) that has all sorts of impact at the level of thermopolitics. Congress might decide to cut programs that fund school meal programs. This, in turn, will have a thermodynamic impact on those students that go without the calories they need developmentally and cognitively to function in a particular way. There is an entanglement here of semiopolitical and thermopolitical domains. The young student here has been constrained both at the level of **semiotic phenomena and thermodynamic structures**. The point is that if true, semiotic intervention (speech-acts, protests, interpretations, deconstructions, etc) will not be an appropriate response to all political problems because social formations are not entirely structured by the semiotic. The child in that school does not suffer from a lack of the right signs, but from a lack of calories needed to run the engine of his thought and body. Certainly semiotic interventions might be needed to render that energy available, but it is the energy itself that is at issue and the absence of that energy that forms the spider web entangling him in his position. A correlationist perspective tends to erase this as even being a site of the political. **Larval Subjects January 28, 2014 Ontological Anarché Posted by larval subjects under Uncategorized January 27, 2014 Roden on Pluralism Over at enemy industry Roden has an excellent post up on the pluralism discussion. Check it out here. January 25, 2014 Different Senses of Pluralism and Ontology Posted by larval subjects under Uncategorized Responding to one of my comments over at his blog, Jeremy Trombley:** I have deleted restricted,

reformulated some sentences .Kindly bear with me. Intention is to make the discursive enucleation consistent with the subject in question**.** Recourse to epistemology of the quantum mechanics generic and fervor of string theory is a sine qua non for the further comprehension of the theory and ontology .How would the world appear to us if its ontology was that of classical mechanics but every agent faced a restriction on how much they could come to know about the classical state? We show that in most respects it would appear to us as quantum. The statistical theory of classical mechanics, which specifies how probability distributions over **phase space evolve** under Hamiltonian evolution and under measurements, is typically called Liouville mechanics, so the theory we explore here is Liouville mechanics with an epistemic restriction. The particular epistemic restriction we posit as our foundational postulate specifies two constraints. The first constraint is a classical analog of Heisenberg's uncertainty principle; the second-order moments of position and momentum defined by the phase-space distribution that characterizes an agent's knowledge are required to satisfy the same constraints as are satisfied by the moments of position and momentum observables for a quantum state. The second constraint is that the distribution should have maximal entropy for the given moments. Starting from this postulate, authors derive the allowed preparations, measurements, and transformations and demonstrate that they are isomorphic to those allowed in Gaussian quantum mechanics and generate the same experimental statistics. They argue that this reconstruction of Gaussian quantum mechanics constitutes additional evidence in favor of a research program wherein quantum states are interpreted as states of incomplete knowledge and that the phenomena that do not arise in Gaussian quantum mechanics provide the best clues for how one might reconstruct the full quantum theory. I: http://dx.doi.org/10.1103/PhysRevA.86.012103 **Reconstruction of Gaussian quantum mechanics from Liouville mechanics with an epistemic restriction Phys. Rev. A 86, 012103 – Published 10 July 2012 Stephen D. Bartlett, Terry Rudolph, and Robert W. Spekkens.** Deleuze in

fact characterizes modes of existence, with their powers and capacities. The answer is this: Deleuze approaches modes of existence, ethically speaking, not in terms of their will, or their conscious decision making power (as in Kant), nor in terms of their interests (as in Marx, for example), but rather in terms of their drives. For Deleuze, conscious will and preconscious interest are both subsequent to our unconscious drives, and it is at the level of the drives that we have to aim our ethical analysis. Here, I would like to focus on two sets of texts on the drives taken, not from Nietzsche and Spinoza, but rather from Nietzsche and Leibniz (Leibniz being one of the first philosophers in the history of philosophy to have developed a theory of the unconscious). The first set of texts comes from Nietzsche's great early book entitled Daybreak, published in July 1881. Nietzsche first approaches the question of the drives by giving us an everyday scenario: "Suppose we were in the market place one day," he writes, "and we noticed someone laughing at us as we went by: this event will signify this or that to us according to whether this or that drive happens at that moment to be at its height in us—and it will be a quite different event according to the kind of person we are. One person will absorb it like a drop of rain, another will shake it from him like an insect, another will try to pick a quarrel, another will examine his clothing to see if there is anything about it that might give rise to laughter, another will be led to reflect on the nature of laughter as such, another will be glad to have involuntarily augmented the amount of cheerfulness and sunshine in the world—and in each case, a drive has gratified itself, whether it be the drive to annoyance, or to combativeness or to reflection or to benevolence. This drive seized the event as its prey. Why precisely this one? Because, thirsty and hungry, it was lying in wait" (D 119). This is the source of Nietzsche's doctrine of perspectivism ("there are no facts, only interpretations"), but what is often overlooked is that, for Nietzsche, it is our drives that interpret the world, that are perspectival—and not our egos, not our conscious opinions. It is not so much that I have a different perspective on the world than you; it is rather that each of us has multiple

perspectives on the world because of the multiplicity of our drives—drives that are often contradictory among themselves. "Within ourselves," Nietzsche writes, "we can be egoistic or altruistic, hard-hearted, magnanimous, just, lenient, insincere, can cause pain or give pleasure" (Parkes, pp. 291-292). We all contain such "a vast confusion of contradictory drives" (WP 259) that we are, as Nietzsche liked to say, multiplicities, and not unities. Moreover, these drives are in a constant struggle or combat with each other: my drive to smoke and get my nicotine rush is in combat with (but also coexistent with) my drive to quit. This is where Nietzsche first developed his concept of the will to power—at the level of the drives. "Every drive is a kind of lust to rule," he writes, "each one has its perspective that it would like to compel all the other drives to accept as a norm" (WP 481). DELEUZE AND THE QUESTION OF DESIRE: TOWARD AN IMMANENT THEORY OF ETHICS Daniel W. Smith. **In on a Possible Physical Metatheory of Consciousness Miroljub Dugic et al** show that the modern quantum mechanics, and particularly the theory of decoherence, allows formulating a sort of a physical metatheory of consciousness. Particularly, the analysis of the necessary conditions for the occurrence of decoherence, along with the hypothesis that consciousness bears (more-or-less) well definable physical origin, leads to a wider physical picture naturally involving consciousness. This can be considered as a sort of a psycho-physical parallelism, but on very wide scales bearing some cosmological relevance. Open Systems & Information Dynamics Volume 9, Number 2, 153 (2002) arXiv: quant-ph/0212128. On a broader perspective and background, the theoretical computation of the universe is justified.IN physics and cosmology, digital physics is a collection of theoretical perspectives based on the premise that the universe is, at heart, describable by information, and is therefore computable. Therefore, according to this theory, the universe can be conceived of as either the output of a deterministic or probabilistic computer program, a vast, digital computation device, or mathematically isomorphic to such a device. Digital physics is grounded in one or more of the following

hypotheses; listed in order of decreasing strength. The universe or reality: is essentially informational (although not every informational ontology needs to be digital); is essentially computable (the pancomputationalist position); can be described digitally; is in essence digital; is itself a computer (pancomputationalism); is the output of a simulated reality exercise. Following Jaynes and Weizsäcker, the physicist John Archibald Wheeler wrote the following: [...] it is not unreasonable to imagine that information sits at the core of physics, just as it sits at the core of a computer. (John Archibald Wheeler 1998: 340) It from bit. Otherwise put, every 'it'—every particle, every field of force, even the space-time continuum itself—derives its function, its meaning, its very existence entirely—even if in some contexts indirectly—from the apparatus-elicited answers to yes-or-no questions, binary choices, bits. 'It from bit' symbolizes the idea that every item of the physical world has at bottom—a very deep bottom, in most instances—an immaterial source and explanation; that which we call reality arises in the last analysis from the posing of yes–no questions and the registering of equipment-evoked responses; in short, that all things physical are information-theoretic in origin and that this is a participatory universe. (John Archibald Wheeler 1990: 5) David Chalmers of the Australian National University summarised Wheeler's views as follows: Wheeler (1990) has suggested that information is fundamental to the physics of the universe. According to this 'it from bit' doctrine, the laws of physics can be cast in terms of information, postulating different states that give rise to different effects without actually saying what those states are. It is only their position in an information space that counts. If so, then information is a natural candidate to also play a role in a fundamental theory of consciousness. We are led to a conception of the world on which information is truly fundamental, and on which it has two basic aspects, corresponding to the physical and the phenomenal features of the world. (Wikipedia). **Participatory consciousness forms the bastion, pillar, post, stylobate and fulcrum of the redressal measures and corrective steps taken for circumvention of**

**many problems, identified and unidentified, spread widest commonalty. Anne Baring's sententious and pithy article will be good to have a timely outlook.** I would like to open my contribution to this meeting with the words of Ken Wilber that I quoted in the first Millennium Symposium: 'The secret impulse of life is towards greater consciousness. Maybe the evolutionary sequence is from matter to body, to mind, to soul, to spirit - each transcending and including, each with a greater depth and greater consciousness and wider embrace' Ken Wilber, and with these from Rupert Sheldrake's contribution to the same meeting: "It is possible to make quite a good case that the sun could be conscious and if the sun, then why not the stars? Why not the galaxies? Why should there not be galactic minds? And if they communicate with each other then the distances are far too great for light to be the medium of transmission. We are into inter-galactic telepathy." After many decades of study, I know of only one diagram in all religious traditions that can adequately explain these statements; that can give us both a blue-print of the many levels of consciousness and a route-map to help us reach beyond our present level of understanding. I would like to introduce a diagram of the four worlds of Kabbalah to illustrate the theme that spirit has brought these four worlds into being and that we, living in the fourth and manifest world know nothing of the existence of the other three invisible ones. It is the clearest diagram I know of that can explain to us the multi-levelled structure of consciousness and our place in the great chain of being. We are indeed, as Wilbur suggests, evolving from matter to body, to mind, to soul, to spirit but where have we come from? It is possible that, as participants in the emanation of spirit into manifest life, we have also come from spirit to soul, to mind, to body, to matter. Everything, from the immense galactic energies to the minute forms of matter participates in one life, is generated from one creative spirit. There is nothing outside this life, this spirit. This diagram suggests that truth is not an article of faith but the experience of participating in the undiscovered reality of what has brought us into being. We cannot know truth until we have entered into

communion with that reality. Continuing in similar vein **Michel Bitbol dilates upon the dialectic deliberation and polemical argumentation between** hermeneutists and eliminativists. When he formulated the program of Neurophenomenology, Francisco Varela suggested a balanced methodological dissolution of the "hard problem" of consciousness. He shows that his dissolution is a paradigm which imposes itself onto seemingly opposite views, including materialist approaches. I also point out that Varela's revolutionary epistemological ideas are gaining wider acceptance as a side effect of a recent controversy between hermeneutists and eliminativists. Finally, he emphasizes a structural parallel between the science of consciousness and the distinctive features of quantum mechanics. This parallel, together with the former convergences, point towards the common origin of the main puzzles of both quantum mechanics and the philosophy of mind: neglect of the constitutive blindspot of objective knowledge. Phenomenology and the Cognitive Sciences Phenomenology and the Cognitive Sciences **Phenomenology and the Cognitive Sciences 2002, Volume 1, Issue 2, pp 181-224 Science as if situation mattered Michel Bitbol**

**ABSOLUTE STABILITY**

Hence if $X_8 / M_8^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_8$ ) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8 - 8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2 - 2y_8\frac{(E_8 + n_8^2 y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right) - 4y_8(E_8 - n_8)$$

And has the solution

$$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8-8)}\left[y_8\left(E_8 + n_8^2\, y_8\right) + \left[y_8^2\,(E_8 + n_8^2\, y_8)^2 + 4\, y_8 (E_8 - 8)(E_8 - n_8)\right]^{8/2}\right]$$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2\, E_8\left[y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)\right]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSOILIDATION**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + \left(M_8 + U_8\right)\lambda_8 - \left(M_8 + V_8\right)\lambda_8 e^{2T_8\lambda_8} + \left(M_8 U_8 + Y_8\right)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8 e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + \left(M_8 + V_8\right) Z_8 \sin\left(Z_8 T_8\right) - \left(U_8 M_8 + Y_8\right) + (M_8\left(U_8 + V_8\right) - X_8)$$

$$\mathrm{X}\cos\left(Z_8 T_8\right) - 2M_8 V_8 cos^2\left(Z_8 T_8\right) + M_8 V_8 = 0$$

$$(M_8 + U_8)\, Z_8 - (M_8 + V_8)\, Z_8\, cos\left(Z_8 T_8\right) + \left(M_8\left(U_8 + V_8\right) - X\right)$$

$$\mathrm{X}\sin\left(\left(Z_8 T_8\right) - M_8 V_8 \sin\left(Z_8 T_8\right)\right) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 =$

$1 - cos\, Z_8\, T_8$and $A_8 = U_8 - V_8$we obtain

$$4 M_8^2 V_8^2 W_8^4 + 2 M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8\left(M_8 + V_8\right) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2 [M_8 - V_8)^2 (2 A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8 (M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8 (A_8^2 + 4(X_8 + Y_8))]$$

$$+ W_8 [A_8^2 + 2X_8)(2 A_8 M_8 - X_8) + 2A_8 Y_8 (M_8 + V_8) - M_8 A_8^3] + A_8^2 (X_8 M_8 - Y_8) = 0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8 > 0$$

$$and$$

$$C_8 > (X_8 - Y_8)/2 M_8$$

$recalling\ that\ A_8 = U_8$ and defining $C_8 = U_8 + V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8 (0, U_8 V_8 = A_8^2 (X_8 + Y_8)$$

$$Q_8 (2, U_8 V_8) = (C_8 + 2M_8^2)^2 (2C_8 M_8 + Y_8 - X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$the existence of a D-position at this point is

determined by letting cos $(Z_8 T_8) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all $T$, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**ABSOLUTE STABILITY IN PREDATOR- PREY MODELS**

**APPLIED TO ACCETUATION ATTRITION MODELS**

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a

marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\,(R) - NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - D(R) - NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The Percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the

decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

## DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t-T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\left(R(t)\right)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The

regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the

**TIME DELAY**

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a dominant eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

DELAY DUE TO CONGREGATIONAL AND CONFEDERATIONAL DEBIT CREDIT TYEPES

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t-T_8)F_8(R_8(t-T_8) - N_8(t)M_8((R_8(t) \quad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8} + \frac{N_8 dG_8}{dR_8} - \frac{dB_8}{dR_8} > \frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8} + \frac{dM_8}{dR_8}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_8 = \frac{N_8 G_8}{2M_8}\frac{d(8nF_8)}{dR_8} (\geq 0) \quad (8a)$$

$$I_8 = -\frac{N_8 G_8}{2M_8}\frac{d(8nM_8}{dR_8} (\geq 0) \quad (8b)$$

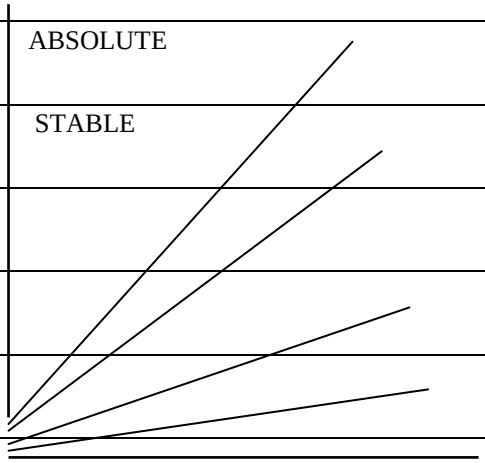
 

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8 > L_8 - I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_8 > L_8 - I_8, \quad \text{when} \quad P_8^2/4 > I \tag{8a}$$

$$P_8 > L_8 - I_8 + \frac{(I_8 - P_8^2}{(L_8 - I)} \quad \text{when} \quad P_8^2/4 < I. \tag{8b}$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit

(like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

See the diagram Nunney gives for predator and prey analysis. We can simulate the same for the financial system in the foregoing too.

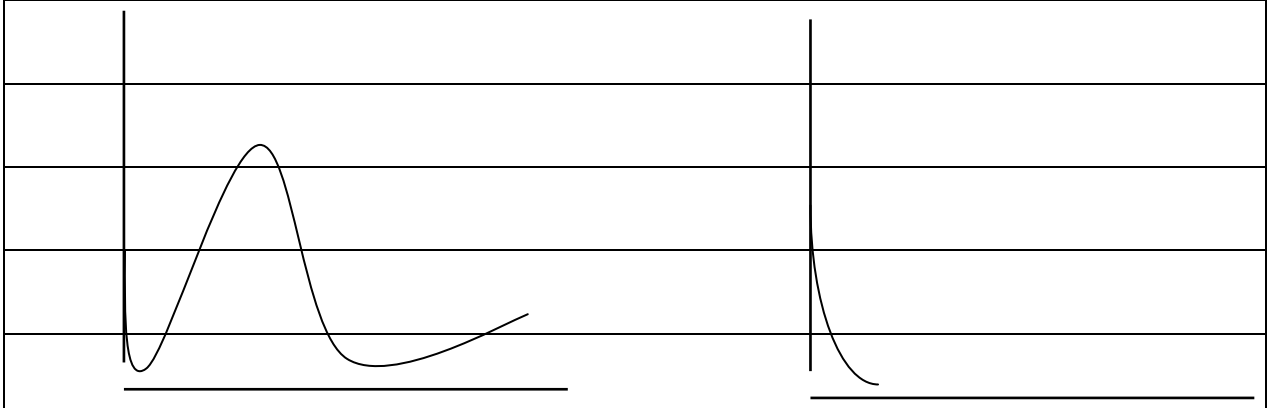
Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the

debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

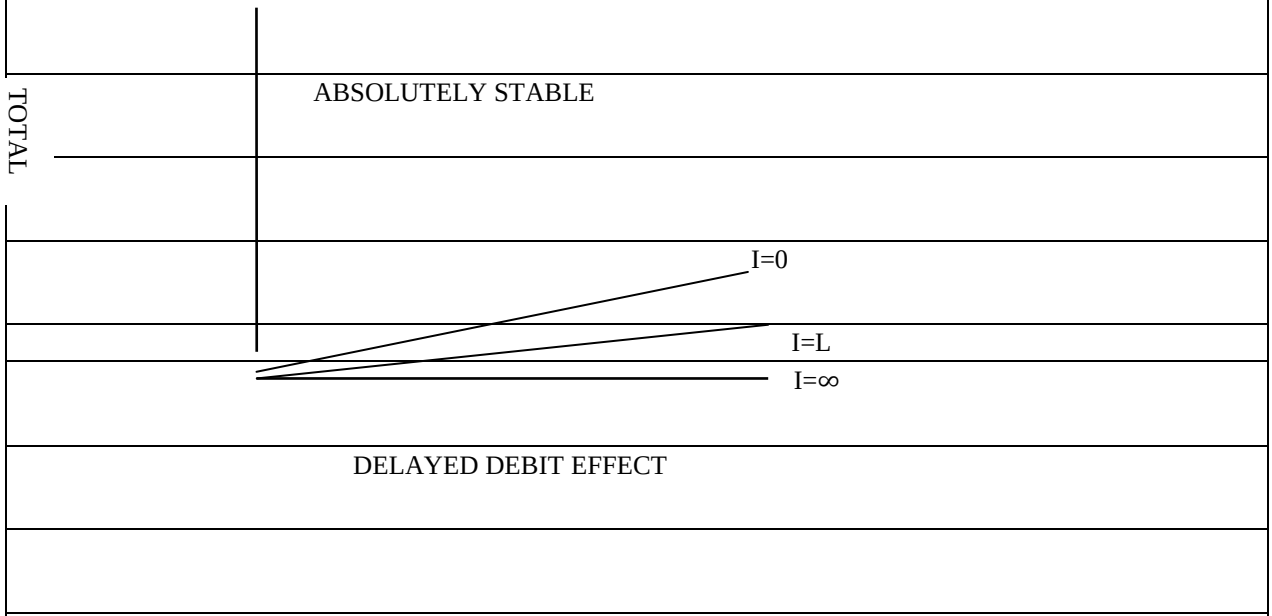
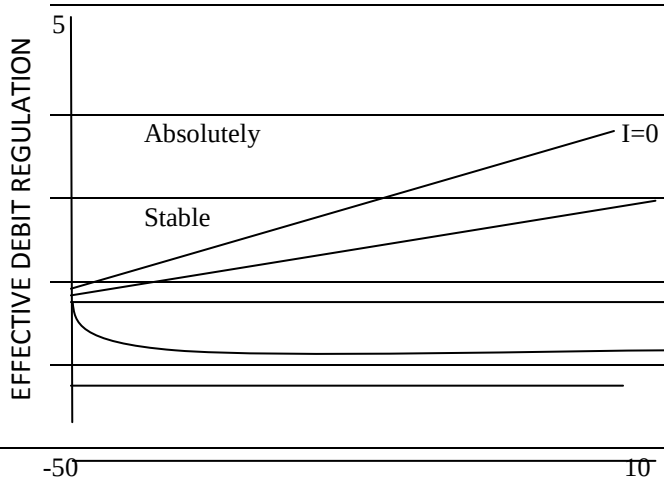



FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison?

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are

defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^\infty \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$ then given the constraints of $f(T_8)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8 \lambda_8 e^{-T_8 \lambda_8} + [(U_8 - V_8) M_8 + Y_8] - [((U_8 -$$

$$V_8 M_8 -] e^{-T_8 \lambda_8} = 0$$

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$\boldsymbol{Z}_8^2 + M_8 Z_8 \sin(Z_8 T_8) - (A_8 M_8 + Y_8) + (A_8 M_8 + X_8) \cos(Z_8 T_8) = 0$$

$$(A_8 M_8) Z_8 - M_8 Z_8 \cos(Z_8 T_8) + (A_8 M_8 - X_8) \sin(Z_8 T_8) = 0$$

Where $A_8 = U_8 - V_8$ Using (A2.3), substitute for $\boldsymbol{Z}_8^2$ and $Z_8$ in (A2.2):

$W_8^2 (X_8^2 - 3 M_8 A_8 X_8 - M_8^2 (X_8 - Y_8) + 2 M_8^3 A_8 + 2 A_8^2 M_8^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_8 = 1 - \cos(Z_8 T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8 / 2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_8 \le 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2 (X_8 - Y_8)$$

$$Q_8(2, A_8) = (A_8 + 2M_8)^2 (2A_8 M_8 - (X_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2 \left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8 \left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8 \left(2X_8 + Y_8\right)\right)$$

$$+2A_8 \left(X_8 + Y_8\right)$$

$$=2M_8^2 W_8^2 (M_8 + A_8) + M_8 A_8^2 W_8 + (2A_8 M_8 - X_8 + Y_8)(A_8 W_8 + M_8 W_8^2$$

$$+ (2 - W_8)((2M_8 X_8 + M_8 Y_8) W_8 + A_8 (X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $W_8$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in *A* will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range $0 \le W_8 \le 2$.Defining

$Q_8(W_8, A_8) = a_8 W_8^2 + c_8$ *where* $a_8$ $b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \square b_8 \square 0$$

And

$$b_8^2 \square 4a_8c_8$$

$$Case(i): Y_8 = 0$$

$Let\ Y_8 = 0$ And try the substitution A= $X_8 = /2M_8 in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8 /$ $2M_8$) is positive for *W<2. Since* $\mathrm{Q}_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\ X_8$ ,

$Let\ Y_8 =\ y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 +\ n_8 y_8)$ we have

$$W_8^2 y_8\left[X_8 n_8 \frac{(1 + n_8 y_8)}{(2 + n_8 y_8)} + M_8^2\left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8 X_8 \left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2}+2y_8\left(n_8-1\right)+\frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)}=0.\ (A2.12)\right.$$

Since $n_8$ and $A_8$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8=(2+n_8 y_8),$ *the condition*(A2.88?)*can be written as*

$[\frac{X_8(E_8-8)}{E_8^2}]^2[\frac{X_8^2}{M_8^2}-4y_8 E_8^2\frac{X_8}{M_8^2}+4y_8(n_8-2)E_8^3]$ ?0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8/W_8^2)$ can be expressed as

$2E_8[y_8E_8+(y_8 E_8(8+y_8)(E_8-n_8))^{8/2}]>\frac{X_8}{M_8^2}>2E_8[y_8E_8-y_8E_8(8+y_8)(E_8-n_8))^{8/2}]$

**ABSOLUTE STABILITY**

Hence if $X_8/M_8^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0?$W_8$ ?2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8-8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2-2y_8\frac{(E_8+n_8^2y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right)-4y_8(E_8-n_8)\ ?0$$

And has the solution

$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8-8)}[y_8(E_8 + n_8^2 y_8) + [y_8^2 (E_8 + n_8^2 y_8)^2 + 4 y_8 (E_8 - 8)(E_8 - n_8)]^{8/2}]$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2 E_8 [y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + (M_8 + U_8)\lambda_8 - (M_8 + V_8)\lambda_8 e^{2T_8\lambda_8} + (M_8 U_8 + Y_8)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8\, e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + (M_8 + V_8) Z_8 \sin(Z_8 T_8) - (U_8 M_8 + Y_8) + (M_8(U_8 + V_8) - X_8)$$

$$\mathrm{X}\cos(Z_8 T_8) - 2M_8 V_8 \cos^2(Z_8 T_8) + M_8 V_8 = 0$$

$$(M_8 + U_8) Z_8 - (M_8 + V_8) Z_8 \cos(Z_8 T_8) + (M_8(U_8 + V_8) - X)$$

$$\mathrm{X}\sin((Z_8 T_8) - M_8 V_8 \sin(Z_8 T_8) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 = 1 - \cos Z_8 T_8$and $A_8 = U_8 - V_8$we obtain

$$4M_8^2 V_8^2 W_8^4 + 2M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8(M_8 + V_8) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2[M_8 - V_8)^2(2A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8(M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8(A_8^2 + 4(X_8 + Y_8))]$$

$$+ W_8[A_8^2 + 2X_8)(2A_8 M_8 - X_8) + 2A_8 Y_8(M_8 + V_8) - M_8 A_8^3] + A_8^2(X_8 M_8 - Y_8) = 0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8 > 0$$

$$and$$

$$C_8 > (X_8 - Y_8)/2M_8$$

$recalling\ that\ A_8 = U_8$ and defining $C_8 = U_8 + V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8(0, U_8 V_8 = A_8^2 (X_8 + Y_8)$$

$$Q_8\left(2, U_8 V_8\right) = \left(C_8 + 2M_8^2\right)^2 (2C_8 M_8 + Y_8 - X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$the existence of a D-position at this point is determined by letting cos $(Z_8 T_8) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all *T*, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**ABSOLUTE STABILITY IN PREDATOR- PREY MODELS**

**APPLIED TO ACCENTUATION ATTRITION MODELS**

We apply model to each and every module and a concatenated model was to be presented which we procrastinate due to spatial restrictions.

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - D(R) - NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

The question is political-discursive. It's what Stengers and Latour are getting at with their concepts of diplomacy and cosmopolitics. They grant, first, that all entities exist and, second, that to say that someone's cherished idol (or whatever disputed entity they hold dear) is non-existent is a **'declaration of war'** – 'this means war,' as Stengers often says. They thus shunt onto-political discourse off of the terrain of knowledge/belief in the sense of existence/non-existence. Their basic claim seems to be that 'respect for otherness,' i.e. political pluralism, can only come from granting the entities that others hold dear an ontology, even if you don't 'believe' in them. You

are thus permitted to say 'I do not follow that god, he has no hold over me' but you are not permitted to say 'your god is an inane, infantile, non-existent fantasy, grow up.' And it's not just a question of politeness (although there's that too). The point is to grant others' idols and deities an existence – one needn't agree over what that existence entails, over what capacities that entity has or what obligations it impresses upon you as someone in its partial presence but to deny it existence entirely is to 'declare war' – to deny the possibility of civil discourse, of pluralistic co-existence. I believe that it was Richard Rorty who once quipped something like claims to reconcile realism and idealism always seem to end with a triumph of idealism. We don't, in fact, get realism through such approaches, but rather just get a pervasive anti-realism. I think this is also the problem with the "non-controversial pluralism" advocated by Stengers and Latour that Phillip defends here. Such pluralism is not realism but is, in fact, a thoroughgoing social constructivism. I think this is the central problem with Latour's argument in Irreductions (these days I regret having ever defended it). In rejecting both Enlightenment critique and what he calls "reduction" he wants to say something like "The Pentacostal really is filled with (e&eb) the Holy Spirit", that for the 19th working scientist heat really is a fluid and phlogiston really is what allows things to burn, and that for the Greek lightning really is an expression of Zeus's anger. Latour tells us that we aren't to reduce or explain away the entities posited by another group's "ontology" but are to develop explanations from within that ontology. It might not sound particularly sexy– and it certainly doesn't tell us what is worth thinking –but I can't help but believe that philosophy is the critical and reflective investigation of basic concepts that guide our investigation of the world about us, how we ought to live our lives, and what form of governance might be best. Compare two figures. A scientist might ask, what causes depression? We can very well imagine a philosopher turning around and asking the scientist, what is causality? The scientist presupposes a concept of causality in her investigations. She uses this concept in her inquiry. Now she might have a sophisticated concept of causality or she might never have thought much about causality at all, using it in the sort of colloquial and unreflective way that Plato decried when, for example, people

like Euthyphro talked about piety. A whole cascade of questions arise when we raise a question like "what is causality?" We can ask whether or not causality exists at all. We can ask how we distinguish between correlation and two events that merely accompany one another from genuine causation. This, for example, was Hume's question. But perhaps most importantly we can ask whether there is only one form of causality or many forms of causality. Is there only one-to-one causation; one cause and one effect? Is there many-to-one causation; or many events conspiring to produce an effect? Is there one-to-many causation; or one event producing a variety of different effects? We can even ask whether causality necessarily moves from past to present or whether there aren't forms of causality that move from future to past! **January 23, 2014 Pluralism and Realism:** multiple-worlds Over at Struggles Forever, Jeremy Trombley has an interesting post up on "the ontological turn" in anthropology or ethnography. I've been meaning to have a discussion with him about this as I think it's an issue many of us are struggling with. For example, the core project of The Democracy of Objects– a project which I think many have missed –is to somehow reconcile some version of social constructivism with a realist ontology capable of making room for ecology (which requires realist and materialist positions as there's a fact of the matter where global warming is concerned) as well as the role played by objective agencies in social assemblages such as technologies, infrastructure, features of geography, local climates, the growth cycles of plants and animals, waste, etc. Maybe we can try to organize some cross-blog event to discuss these issues. I certainly think they're close to the heart of Jeremy, Michael of Archive Fire, Arran James, and a host of others. As an aside, I'm beginning to realize how the different sites of the political I've been outlining– semiopolitics, thermopolitics, oikopolitics (political economy), geopolitics, eropolitics (the politics of sex and desire), biopolitics, and chronopolitics (and I'm sure there are other political sites!) –are drawing me away from traditional Marxism. Assuming that classical Marxism holds that economics or the conditions and relations of production are determinative of all other sites of the political, the various sites of the political that I've been outlining would lead to the conclusion that there is not one determinative

base of the political. This would not require committing Marx to flames, but rather of recognizing the phenomenon of overdetermination, or of a variety of different entangled sites of the political. But I digress. First, I find myself wondering what the ontological turn means in ethnography. Is it 1) the investigation of the different ontologies held/proposed by different cultures? E.g., the Aztecs believed that reality was structured in this way, while the Greeks in that way, and the ancient Chinese this way, etc? Or 2) is it an investigation of how real entities– independent of cultural beliefs –influence cultural formations? Or is it a combination of both? A position that I would favor. January 20, 2014 Model is applied to structured reality

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

**The evaporation spectrum of black holes from a local quantum gravity perspective Aurélien Barrau (Submitted on 21 May 2016**)

(1) Authors revisit the hypothesis of (e) a possible line structure in the Hawking evaporation spectrum of black holes, due to (e) non-perturbative quantum gravity effects, even arbitrarily far away from (e) the Planck mass.

(2) Authors show that this naive prediction might in fact hold in (eb)the specific context of loop quantum gravity, with (e&eb) a small departure from (e) the ideal case for (e) some low-spin transitions.

They also show that the effect is (=) neither washed out by the dynamics of (e) the process, nor by (e) existence of a mass spectrum up to (e) a given width, nor by the secondary component induced by (e) the decay of neutral pions emitted during (e&eb) the time-integrated evaporation. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Cite as: arXiv:1605.06613 [gr-qc] (or arXiv:1605.06613v1 [gr-qc] for this version

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t - T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\left(R(t)\right)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

It's very easy to treat this as an abstract, academic debate: "Are you a realist or are you an anti-realist?", as if it were just a matter of what happens to be true. But it seems to me that this debate has, in reality, always been about politics. As I outlined in my last post, we have perpetually seen how appeals to the real and natural have been used in the name of oppressive power, inscribing both the exploitation of nature and the oppression of various people in the very fabric of being itself. Theistic theology and realist ontology have perpetually been used in the name of

what Deleuze called "State Philosophies" or philosophies that ontologize contingent orders of power and privilege (e.g., "the great chain of being" used to justify patriarchy, monarchy, serfdom, poverty, etc, and appeals to nature used to justify poverty and racial inequality (The Bell Curve), patriarchy (evolutionary psychology), heteronormativity, capitalism, etc). Because arrangements of power and inequality are always contingent in the sense that there's no marked difference in the capacities of peoples, power always looks for a transcendent supplement that would provide justification through ontological necessity. Antirealism– from the Greek atomists to present –became the radical and emancipatory gesture because it revealed the lie behind all of these forms of social organization or their inherit contingency or arbitrariness. Realism, by contrast, has all too often functioned as an apologetics for arbitrary power and social organizations. Here it's worth recalling what Foucault said about science: "…Even before we know to what extent something like Marxism or psychoanalysis is analogous to a scientific practice in its day-to-day operations, in its rules of construction, in the concepts it uses, we should be asking the question, asking ourselves about the aspiration to power that is inherent in the claim to being a science. The question or questions that have to be asked are: "What types of knowledge are you trying to disqualify when you say that you are a science? What speaking subject, what discursive subject, what subject of experience and knowledge are you trying to minorize when you say 'I speak this discourse, I am speaking a scientific discourse, and I am a scientist.' What theoretico-political vanguard are you trying to put on the throne in order to detach it from the massive, circulating, and discontinuous forms that knowledge can take?" (Society Must Be Defended, 10). All of these questions hold equally for claims to something being real. What is one trying to minorize when claiming something is real? What becomes privileged? What is excluded? It is these questions that have been at the heart of the realism debates, for as Spencer-Brown taught us, every distinction has a marked and unmarked space, draws attention to something to be included and pushes something into the unconscious or the domain of the invisible, hidden, or veiled. This is above all the case with evocations of the real. However, as I've tried to show antirealism

leads to its own problems. First, so long as we exclude real beings from our ontological inventory, we are unable to fully understand how power functions ("such and such a set of cultural formations have been thoroughly debunked, yet people still live as if they believed them"). Not only do we not fully understand the sources of the problems due to too much focus on the discursive and semiotic, but we deny ourselves valuable sites of political intervention at the level of infrastructure. Second, we are prevented from addressing things such as cultural racism, such as that found in Heidegger's privileging of the West and the Greeks and Germans in particular. We do a good job addressing biological and theological racism, heteronormativity, and sexism by showing how it is a cultural construction, but when faced with racism such as Heidegger's where he argues that there's something "unique" about the Greek event, it's language, and about the German language, or Badiou/Zizek's racism with respect to the "Pauline Event", we really have no response. Here someone like Jared Diamond or Fernand Braud el is needed to explain global-geographical inequalities. Third, the tools of the cultural turn really do not provide us with the means of thinking the ecological as a site of the political. For this reason, I've tried to formulate a third way– which might be called "constructivist realism" or "constructivist naturalism" –that retains the insights of the cultural turn, while also allowing a robust place for the material. I don't claim to be original in this. I think that many such as Manuel DeLanda, Deleuze and Guattari, Stacy Alaimo, Karen Barad, Jane Bennett, etc., are up to something similar. It's a vast project that requires the work of a multitude of voices, especially given the way in which the culturalism pervades humanities. **The Relational and the Non-Relational: Notes towards an Immanent and Pluralist Theory of Meaning Posted by larval subjects under uncategorized k-bigpic I have deleted or reformulated some sentences for which I beg pardon attributable to spatial constraints** What is correlation, what is causation? What is coordination? What is coincidence? These thoughts have bothered humans since long. I have been able to excavate some reliable material from Wikipedia which I would like to share. "All views are partial and contingent – that's the lesson of pluralism, and you say as much yourself when you say that we must be attentive to ways our own

knowledge might contain superstitions. But it's not discourse that determines whether one view works and another doesn't – it's a confrontation with **non-discursive (non-human)** agencies." I wonder if part of the issue here is that we understand pluralism differently. For me this doesn't sound like pluralism at all. Rather, it just sounds like our **epistemological condition**. We [hopefully] come to understand those portions of existence we question or investigate. The reason we come to investigate them is largely contingent. Finally, the accounts of these features of existence we give can be mistaken. All of this is perfectly consistent with a monism. I'm not sure why what he outlines above is the lesson of pluralism. It seems to me that every realist knows this. It seems to me that ontological pluralism is something quite different. Ontological pluralism is the thesis that there are many different worlds inhabited by many different entities. Thus, for example, you would have one world; say that of Lucretius, that's only inhabited by atoms and their combinations. You would have another world; say the world of the Mongolian shaman, that's inhabited by spirits that do all sorts of things. The ontological pluralist is saying that these spirits are, that they exist, that they're real. This is something quite different than merely saying that there are partial and contingent points of view on the world. It's this saying that these entities are rather than that some person or groups of people believe that they are that is the nub of the issue. Now, I think part of the issue here is that there's an ambiguity in the term ontology. Ontology can be one of two things. On the one hand, ontology is a group or persons set of beliefs as to what is. Here it's trivially true that there are pluralities of ontologies and the realist readily recognizes this. This is the whole reason there are debates over ontology. Mongolian shamans have their ontology, Europeans theirs, Christian fundamentalists theirs, materialists theirs, etc. When striving to understand and communicate with others it's vital to understand these ontologies because, as rhetoricians like Burke point out, our **beliefs about what is** are among the things that motivate our action. Despite having never seen bacteria I was my hands because I believe there are bacteria and viruses on door handles and whatnot and don't want to get sick. My belief about a particular thing existing is what motivates my action. And who knows, perhaps this belief is as

superstitious as the belief that the crops failed because God was displeased with my community. On the other hand, ontology is a theory **about what is.** It is making a claim that something exists. This is where the rubber hits the road. The ontological pluralist seems committed to the thesis that every group set of beliefs about what exists is sufficient for granting the existence of those entities. This is what I find objectionable in Latour. Obviously I'm not bothered by Latour's suggestion that we should take into account the role that nonhumans like speed bumps, rivers, microbes, etc., play in the form that social assemblages take. But this is not all that Latour claims. This is precisely where the philosopher might balk. Unlike the ethnographer, the philosopher is not interested in what people believe exists, but rather philosophers– at least of the realist variant –are trying to figure out what is. In other words, the realist philosopher begins with the premise that not all of these beliefs about what is are true. So for the philosopher, recognizing that Mongolian shaman's believe in the existence of shamans would only be the first step. The next step would consist in determining whether there's good reason for thinking such entities really do exist, i.e., whether there's good reason for believing these entities have mind and culture independent reality. Lest readers think that I'm just picking on the supernatural here, we can ask similar questions about strings, subatomic particles, galaxies, etc., etc., etc. I guess the question really comes down to what exactly we mean by ontological pluralism. When we talk about ontological pluralism are we defending the thesis that people have different theories of being? If so, then ontological pluralism is trivially true. If this is what is meant, then I certainly share Jeremy's view that it's valuable to understand the different world's people believe in. Certainly when I was practicing as an analyst I didn't get in ontological debates with my patients and it was necessary to understand their theory of being or their ontology to properly attend to them. Or, when we talk about ontological pluralism are we defending the thesis that all these ontological theories are true and refer to really existing entities? That's quite a different claim and is not one I would defend or endorse. Now someone might object that "in both The Democracy of Objects and Onto-Cartography you defend the thesis that there are multiple worlds." This is true. Because I hold that not everything

is related I'm led to the conclusion that there are diverse worlds. However, I also hold that however many worlds there might be, these worlds are nonetheless composed solely of material entities. Within the framework I propose I wouldn't suggest that there's one world where there are spirits and another world where there are souls and yet another where there are only material entities. My view is that there aren't spirits or souls in any of these worlds. One of the things I keep hearing in these discussions is that somehow the realist adopts a view from nowhere. I honestly don't understand this criticism. Investigation always occurs somewhere and requires all sorts of mediations involving technologies, experiments, etc. It's that labor of gathering evidence, conducting experiments, using technologies to observe the world, etc., that gradually gives us a body of data that allows us to say there's good reason to believe that such and such a thing exists and has these powers. Another charge seems to be that the realist refuses to recognize that their claims about the world are fallible. I find this charge particularly strange because it's precisely because the realist recognizes the difference between our theories of the world or what we say about the world and the world itself that fallibility is built into the core of his position. Realism doesn't mean one holds they have special access to the world, that they know all truths, or that they have the truth in hand, only that there are truths to be known and that we can be mistaken about things. Circling **Squares Philip** has a post responding to my quandaries about how to mesh realism and pluralism. He writes: Ontologically and metaphysically the idea of realist pluralism is no longer an issue. There are (appropriately) numerous variants but the basic idea that reality is itself pluralistic is well established. Captioned expatiation is about Everett worlds and string theory

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from

being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

The von Neumann formulation leads us to conclude that the lower realms of quantum reality form a virtually limitless array of potentials that are inherently incapable of realization without the observation of a knowing entity, which we identify with consciousness (Kafatos and Nadeau, 1990). We thus have a correspondence between quantum theory, the hierarchical structure implied by the Holographic Principle, and the higher orders of experience,

leading to the full expression of the Conscious Universe. David Bohm's theory of the implicate and explicate orders involves a holographic Principle that is fully consistent with the Holographic Principle discussed here (Bohm 1980; 1986; Bohm and Hiley, 1993). According to Bohm, there is an implicate order that represents the universal, holographic subtext of reality, and which unfolds in every moment to produce the explicate order that we all observe. Thomas Germinario (2004) has equated the implicate order with (eb) the unconscious process, and the explicate order with (eb) conscious process. He emphasizes the nature of dreams within (eb) the implicate order, and the importance of the dream work for maintaining a healthy mind through (e&eb) integration of the implicate order and our daily lives. Allan Combs and Mark Holland (1990) connected the implicate order or holomovement (Bohm, 1980) with (e&eb) Carl Jung is theory of synchronicity (Jung and Pauli, 1955), with (e&eb) the implicate order providing (eb) a holographic medium through which apparently disconnected individuals become connected. The principle of synchronicity, the instantaneous connection of people and events beyond (e) the senses, has been equated with the quantum-physical principle of non-locality (Combs and Holland, 1990; Germine, 1991), and has been proposed to be the fundamental mechanism of conscious process (Germine, 1991).Applying the Holographic Principle theory of mind, subjects may be connected synchronously in the manner of two individuals conversing through cellular phones, where the radio signals are transmitted through (e&eb) a distant satellite. The difference here is that the radio signals travel at the speed of light, and thus the there is some miniscule time lapse between the sending and receiving ends. The Holographic Principle implies (eb) a much more distant and instantaneous communication pathway, with (e&eb) the analogue of the satellite being the holographic boundary of the Universe, as well as a much richer communication, potentially involving a transmission of thoughts and feelings. Our ego-consciousness seems to mask (e) the universal relatedness implied by (e) the Holographic Principle, and it is perhaps only through transcendence of (e) the ego-consciousness that the higher orders of experience can become conscious. In the phenomenon of synchronicity, there seems to be a meaningful connection between (e&eb) individuals that

breaks through the barrier of ego-consciousness. Such a connection is reported by many individuals in the course of dreams, when (e) the ego-consciousness has been suspended, at times informing (e&eb) the dreamer of something that has happened in the life of a meaningfully-connected individual. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Impassioned-intense Connectionist neuroscience, in an effort to attain relevance to clinical practice, has recently given us a theory of the relationship of the unconscious to (e&eb) consciousness which is (=) dualistic (Viamontes and Beitman, 2007), an ad hoc theory of conclusions based on (e) assumptions, the most erroneous of which are that representations of (e) information are information, and that such representations elaborate themselves through (e&eb) a process of representations of representations in neuronal circuits (Bechtel and Abrahamsen,1991). As theory, the connectionist models are lacking (e) heuristic value and **are (=) inimical to the progress** of (e) basic and clinical neuroscience. This is not to denigrate or minimize the importance of neuronal connections, but these connections are rather like the wires in (eb) a radio that plays a symphony. The radio, in and of itself, is (=) incapable of playing (e&eb) the symphony. The radio doesn't do not write the symphony, nor does it broadcast the Passionate-raging symphony over the airways. Yet, if we loosen one wire, the symphony is (=) not manifested through the radio. The brain is (=) not subject to (e) the same vulnerability, in that it has (e) parallel and redundant networks of processing, and, it is in this sense that the connectivity of the brain becomes (eb) important. Connectionism, to the extent that it accepts (e) consciousness and the (e&eb) unconscious at all, attributes them to (e) separate groups of "circuits" in the brain, with (e&eb) consciousness-processing circuits having (e) a limited but detailed ability to analyze (e&eb) information, and unconscious-processing processing more information in less detail (Viamontes and Beitman, 2007). There is an element of truth in this view in that information that has importance and meaning resonates within the brain at higher levels of experience and at higher levels of holographic recursion. This resonance involves what we know, what we feel, and what we are capable and willing to consciously realize. The Universe, as William Blake noted,

can be seen in a grain of sand. Indeed, the grain of sand requires (e) the Universe to exist. The human organism requires (e) the Self to exist. The Self is not a simple amalgam of (e) sensory experiences, as suggested in the connectionist model (Blinder, 2007). This amalgam view of the Self can be damaging when applied to (e&eb) psychotherapy, as the connectionists suggest (Blinder, 2007). The Self is ours forever and ever, from eternity to eternity. As Milton suggested, the mind is (=) its own place. Our deepest purpose is for our minds to resonate with (e&eb) the supra-conscious levels of experience, much as it is (=) purpose of the radio to play the symphony. If the symphony is not to our liking, (eb) we can turn the radio off. So, perhaps, the transpersonal consciousness can turn (e&eb) our own individual and collective radios off, or perhaps turn the volume down, if (e) we are not in harmony with the **Universal Mind**. Physics has given us many enigmas. Quantum theory itself is an enigma. But some of the enigmas seem to involve (e&eb) the limitation of (e) our own abilities of discernment. The thought experiment of (e) "Schrödinger's Cat," as it is often interpreted (e.g. Gribbin, 1984), assumes (eb) that the **cat is not conscious, and is (=) incapable of observation in the context of (e) collapse of (e) the wave function.** This is due, in large part, to the false dichotomy of (e) information and (e&eb) experience. Our observation of other animals indicates, to those that are attuned to (e) the animal mind, who can "feel" the animal mind, that, at (eb) the very least, all mammals are (=) conscious. It would seem that the great neuropsychiatrist, Stanley Cobb (1948), may have been correct in his attribution of (e) levels of consciousness to a wide variety of organisms: "...lower animals with (e&eb) no cerebrum appear to be conscious...even plants such as the sunflower that turns towards strong light may have (e) a vague awareness or warmth and comfort. There are many degrees of consciousness and it is my contention that it is integrated at (e&eb) many levels like other important functions of (e) the central nervous system."With respect to the theory of mind, recent mainstream thinking remains (eb) classical and mechanistic. Such views of mind are advocated by our most in prominent neuroscientists (e.g. Changeux, 1985; Gazzaniga, 1985) as neuroscience research attempts to explain all mental function on the basis of (e) **pure mechanism and the localized function of**

**specific areas of the brain.** Francis Crick, the distinguished Nobel Laureate, had a very influential second career as a neuroscientist. He concluded that we are basically soulless creatures, that all mental processes could be reduced to (e&eb) identifiable, neural correlates in (eb) the brain, and that, in fact, we are (=) nothing but a collection of neurons (Crick, 1995). Crick even went so far as to say that what we call the soul is (=) a group of neurons located in the prefrontal lobes of the brain. Another Nobel Laureate, Gerald Edelman (1987), has written a critically acclaimed book on "neural Darwinism," which purports that the brain is "circuitry" is constructed through a developmental process of "survival of the fittest" neurons and neuronal connections. At the same time the total purposelessness and lack of direction in evolution is expounded, again to critical acclaim, by Richard Dawkins (1986). The idea that there essentially no self or soul imbued with agency and self determination was expounded by Gilbert Ryle (1949/2002) with his parody of the "ghost in the machine," and this work continues to be influential among many scholars of the mind. Philosophers such as Patricia Churchland (1986), who coined the term "neurophilosophy," take an "eliminative materialist" view of (e) belief, free will, and consciousness. Other philosophers, such as Daniel Dennett (1991) would make (eb) experience "epiphenomenal," a by-product of (e) brain processes with (e&eb) no effect whatsoever on (eb) the function of mind. However, unlike neuroscience, which is still in its infancy, the philosophy of mind has had a long history, including such notable figures as Plato, Aristotle, Descartes, Kant, Locke, Hume, Kierkegaard, and James, to name just a few. Perhaps, as noted by Ervin Laszlo (1974): "in today's world, most of the traditional functions of (e) cognitive synthesis have atrophied (e) and are ignored and neglected." Hopefully, what we have presented here will be a step in the right direction for what Laszlo calls a "conceptual synthesis" to help fill the need for (e) meaningful engagement is such a world. The fundamental elements of such a synthesis have been described by (e) Laszlo as follows: Conceptual synthesis performs (e&eb) at least five basic functions in the guidance of affairs. They are the mystical, the cosmological, the sociological, the pedagogical or psychological, and the editorial functions. The mystical function inspires (eb)

in man a sense of mystery and profound meaning related to (e&eb) the universe and of himself in it. The cosmological function forms (eb) images of the universe in accord with (e&eb) local knowledge and experience, enabling (eb) men to describe and identify (e&eb) the structure of (e) the universe and the forces of nature. The sociological function validates supports and enforces (eb) social order, representing (eb) it in accord with (e&eb) the nature of the universe, or as the natural or right form of (e) social organization. The pedagogical or psychological function guides (eb) individuals through stages of life, teaching ways of understanding themselves and others and presenting (eb) desirable responses to life is challenges and trials. Finally, the editorial function of (e) conceptual is to define (eb)some aspects of reality as important and credible and hence to be attended to, and other aspects unworthy of serious attention. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Kant sums up this inversion and its spirit early in the Critique of Pure Reason: Up to now it has been assumed that all our cognition must conform to the objects; but all attempts to find something about them a priori through concepts that would extend our cognition have, on this presupposition, come to nothing. Hence let us once try whether we do not get farther with the problems of metaphysics by assuming that the objects must conform to our cognition, which would agree better with the requested possibility of an a priori cognition of them, which is to establish something about objects before they are given to us. This would be just like the first thoughts of Copernicus, who, when he did not make good progress in the explanation of the celestial motions if he assumed that the entire celestial host revolves around the observer, tried to see if he might not have greater success if he made the observer revolve and left the stars at rest. (B xvi) Onticology, like all variations of object-oriented ontology, is realist in its orientation. In defending a realist ontology onticology holds that the vast majorities of objects, actants, beings, or entities are independent of humans and are what they are regardless of whether any humans regard them or register them. In short, onticology rejects any anthropomorphic, idealist, or anti-realist thesis to the effect that to be is to be the correlate of mind, spirit, the body, the human, and language or otherwise. While it is certainly the case that knowledge is necessarily dependent on the object to which it relates, the converse does not hold true. Objects are not dependent on being known, regarded, perceived, or spoken

about. As such, and to put it in Aristotlean terms, knowledge is an accident of objects, not objects an accident of knowledge. As Althusser so nicely puts it, "[n]o doubt there is a relation between thought-about-the-real and thisreal, but it is a relation of knowledge, a relation of adequacy or inadequacy of knowledge, not a real relation, meaning by this a relation inscribed in that real of which the thought is the (adequate or inadequate) knowledge" (Reading Capital, 96). Althusser goes on to remark that "[t]he distinction between a relation of knowledge and a relation of the real is a fundamental one: if we did not respect it we should fall irreversibly into either speculative or empiricist idealism" (ibid.). Onticology categorically endorses Althusser's verdict. It is a fundamental necessity to distinguish between those relations that belong to the object and those that belong to knowledge. Contemporary philosophy continuously confuses these two very different sorts of relations. Naturally the question arises of how it is possible to surmount our relation to the object so as to determine whether objects themselves possess the properties we encounter in relating to objects. In other words, given that we can only ever relate to the object in relating to the object how is it possible to surmount this relation to get at the being of the object itself? Much more will have to be said about this later– and the answers will be surprising with respect to standard prejudices about realism –however, for the moment it can be said that onticology takes its epistemological inspiration from the transcendental realism of Roy Bhaskar. Among other things, Bhaskar sought to provide a transcendental grounding for the sciences. Insofar as onticology defends the thesis that the field of being is much more vast than the field of objects investigated by the natural sciences, it parts way with the thesis that the domain of being is exhausted by the domain of natural objects. However, the general form of Bhaskar's argument holds for our realist purposes. A transcendental argument seeks to elucidate the conditions under which certain acknowledged practices and forms of cognition are possible. Kant, for example, asked what must be the case for mathematical judgments to be possible. How is it both that we are able to extend our knowledge, as if by magic, through mathematical judgments and, more significantly, that these judgments are able to provide genuine knowledge of the world despite the fact that these

forms of reasoning are not based on experience? Part of Kant's argument consisted in claiming that mind imposes the forms of space and time on the data of experience. In other words, space and time are not attributes of being itself but rather of the mind that regards being. Insofar, Kant argues, as mathematics is ultimately a rumination on the nature of space and time taken in their most abstract form and insofar as the mind imposes space and time on the manifold of sensation, it thus follows that a priori judgments about the nature of spatio-temporal relations are possible that anticipate the structure of actual-space times without directly experiencing these space-times. Why? Because any manifold of sensation must necessarily be structured by these forms imposed by intuition. **Onticology–A Manifesto for Object-Oriented Ontology Part I Posted by larval subjects under Object-Oriented Philosophy My Question Is This: You Say, "Granted There Are Deeper And Deeper And More Primitive Forms Of "Knowledge" Or Abstraction Into Mechanism, But I Wouldn't Say That The Atomic Level Brings A New Epistemic World Into View." But Doesn't Any Kind Of Knowledge, However Primitive, Imply Bringing Some Kind Of Epistemic World Into View Or Proto-View?** Joel: David, Yes Exactly. My Question Is, Just How Far Down The Complexity Gradient Can The Meaning Of 'Knowledge,' Or 'Semiotics, Or Specifically Representation (Maya) Be Stretched Without Breaking? I Can Comfortably Stretch It To The Genetic Level, As The Proto-Epistemic, Where Genetic Codes 'Represent' Phenotypic Structures Or Modes, And The Whole Of Evolution Learns From Its Mistakes And Successes As That Knowledge And History Is Encoded Into And As The Recapitulation Of The Embryogenesis Of The Organism. But An Atom, In My Model, While Certainly Reacting To Its Environment From Its Own Attractor, Essence, And 'Interiority' (Prehension) Isn't Making Representations Of It To Do So. But The Key To The Power And 'Essence' Of Representation (And Simultaneously Its Limitation) Is The Capacity Of Abstraction And Choice. Evolution Does This At A Rudimentary Level (And Anthropomorphizing A Bit) With Its Sending Into The World Its 'Random' Variations, Organismic Ideas, In A Sense. Its 'Mind' Is Literally On The Outside As Our Living World And Biosphere. A 'Choice' Is Made In The Biospheric Interiority,

And Knowledge For The Evolutionary/Embryogenetic Trajectory Is Gained By The "Differential Reproductive Success" Of These Individuals Lives Through Time, As The Increasing Intelligence, Representation, And Capacity For Choice Is Directly Injected Into The Flow Through Sexual Selection. But Where Is The Encoded/Abstracted (Enfolded) Representation Of Options And A Choice Among Them Made At The Atomic Level? I Am Working With The Most Cutting-Edge And Coherent Models For The Atom (Including My Own) That I Can Find Among The Heterodoxy, And I See No Evidence Of Representation And Choice At This Level, Or Rather No Real Way To Further Stretch The Meaning Of The Terms To This Level. Sure There Is Uncertainty From Infinite Difference And Immanent Causation (See Bohm On Infinite Causation), But Without Representation There Really Is No Choice (Although Crudely We Could Call A 'Bifurcation' A 'Choice' By The System As A Whole). So I Am Simply Saying That Representation At The Atomic Level Is Entirely Enfolded And These Are Really Just Ontic Level Phenomena. The Atomic Level, In My View, Is Pre- Not Yet Proto-Epistemic, Which I Place at the Genetic Level Where We Find the First Codes And (e&eb) Primitive Representation. This Just Means That There Is A Real Distinction Between Ontic And Epistemic (And A Gradient Between Them) And We Can't Collapse The Two Into Just The Ontic-Epistemic. And Yet, With The Distinction Intact, And The Gradient To Explore, They Remain Nondual, Univocal, ONE Or Ontic. What I Would Say Exists At All Levels (My Own Self-Aware Myth About The Given) Is What I Call The "Symbiogenesis Of Subject And Object". This Is The Spinozan Essence Of Dynamic Stability And Growth, And Leibnizian Prehension. Also Called The "Nucleation Of Observability" In Spinbitz. It Is From This Nucleus Of The Observer (Fuller) In A "Point-Free Geometry" (Whitehead) Of Pointless "Points Of View" As Conditions Of Boundary That Wilber's "All Is Perspective" Finds Its Ontic Grounding. And This Symbiogenesis Is Also A Rudimentary Or Deep-Level Native Or Primitive Intelligence, Evolution At The Involution Of The Subject-Object Interface. It Is An Exploration Of Creative Learning And Direct Immanent Awareness, Unmediated (Enfolded) By Representational Forms (Unfolded). It Just Is, In Waves (Cosmic Vertebrae)

Before It Abstracts, Represents And Incarnates The Flesh Of What Could Be. Matter Is Enfolded Seed Of The Flower Of Its Abstraction. Maya Is This Depth And Recursion Of Brahma Into And Through The Boundary Conditions Of Abstract Relation. And Brahma Is The Span Of The Emergent Rigor Of The Cosmic Ergodic Spine That Opens The Recursions Into The Field Of Infinite Difference. In Immanent-Transcendent Waves, Literally And Empirically A Limit-Cycle In Ergodicity, It Enfolds And Unfolds From Anatom To Anatomy, From Infinite Intelligence Or Omnirelational Awareness Of Self, Through The Representational Depths Of Self As Self-Consciousness. Omni-Evolution, Trans-Dynamics, And The Anatom Posted On February 21, 2013A Discussion With Tom Huston, Part II The Following Is Part 2 Of An Excerpt From A Discussion Initiated By Tom Huston On His Facebook Page, Edited And Embellished For Clarity. Find Part One, Here. In This Part We Dig Into The Nature Of The "Anatom" And The "Ergodic Spine" Of An Electro-Fractal Cosmos, And Into The Trans-Dynamic Integration Between Being And Becoming, And The X-Interface−The Crossroads Of The Ontic-Epistemic (Brahma And Maya) And The Subject-Object Polarities. Univocal_Dynamics_V1 Tom Huston Is A Founding Member Of Integral Institute And A Former Editor Of Enlighten Next Magazine. His Webpage Is Tomhuston.Com. David Marshall Is A Writer And Editor Living In Chicago. No meaning, value, or normativity to be found in nature, that there is nothing **natural beings ought to be,** but that, rather, these judgments arise from us (See last post of larval subjects) which has led some to raise valuable questions about the coherence of these claims. The problem arises when the following three propositions are taken together: 1. there is nothing outside of nature.2. Beings have no intrinsic meaning, purpose, or value (in and of themselves, there's nothing they ought to be). 3. Value judgments about what beings and being ought to be arise from us and beings like us (bonobo apes, dolphins, institutions, birds of paradise, etc). The problem arises between proposition 1 and 3. How can it both be true that there is nothing outside of being and those normative judgments belong to us and other beings capable of making normative judgments, not nature? The problem arises from restricting these judgments to humans and beings capable of

making these judgments. In making such a claim it seems as if we're saying that there's something outside of nature, something that is beyond nature, thereby violating the first thesis and potentially reintroducing the nature/culture distinction. By contrast, if we say that normative judgments are the special domain of those living beings with the proper degree of sentience to make such judgments, then we seem to reintroduce the nature/culture distinction and fall back into the sorts of problems that I outlined in my last post (and that are so nicely critiqued by thinkers such as Latour). Is there a way out of this? I don't know. Great_Chain_of_Being_2_ (lighter) ASIDE: It seems to me that this is really what the debate between realism and anti-realism, realism and socio-linguistic constructivism surrounding the new materialisms and speculative realists has really been all about.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \qquad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9}+\frac{N_9 dG_9}{dR_9}-\frac{dB_9}{dR_9}>\frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9}+\frac{dM_9}{dR_9}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usance bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_9=\frac{N_9 G_9}{2M_9}\frac{d(9nF_9)}{dR_9}(\geq 0) \qquad (9a)$$

$$I_9=-\frac{N_9 G_9}{2M_9}\frac{d(9nM_9}{dR_9}(\geq 0) \qquad (9b)$$

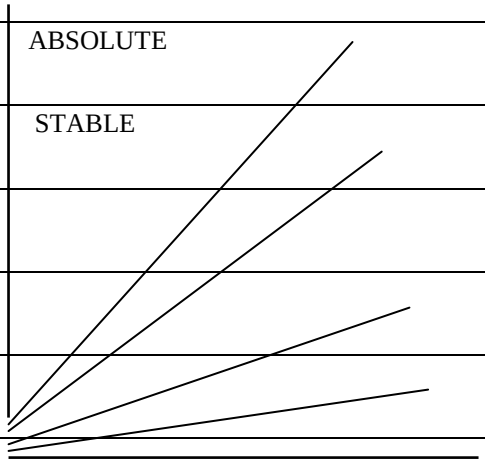
It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9>L_9-I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_9>L_9-I_9$, when $P_9^2/4>I$ (8a)

$P_9>L_9-I_9+\frac{(I_9-P_9^2}{(L_9-I)}$ when $P_9^2/4<I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Reffer couffesy Nunney

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability

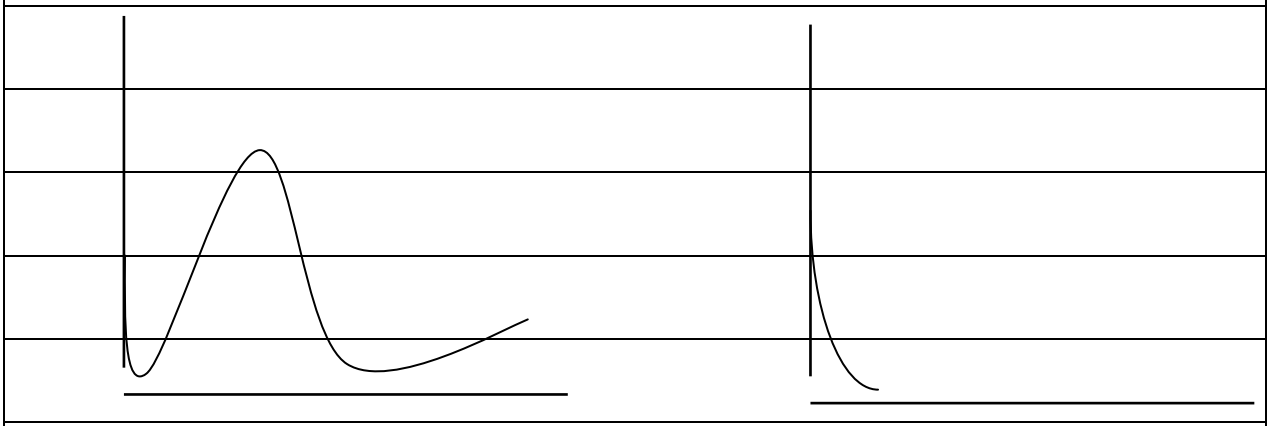
 

is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

Indeed, Nunny's model can be applied to some of the consciousness and philosophical problems. Predator prey analysis has ripped open the veneer of the fundamental concept of universe namely Debit-Credit world.

The noise of ordinary consciousness and (e&eb) beliefs drowns out (e) the signal. Unconsciousness is the background of (e) our ordinary awareness. Our organism is very much at (eb) the center of such effects. The organismic source is (=) our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is a meme-Scape. Pre-conceived concepts vie with structures, concepts with images. Like scientists who ignore assumed truths, we leapfrog over our beliefs and personality deficits, claiming idiosyncratic imagination is literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is both. We inherit the structures that make our experience what it is.

But the structure itself is “empty,” and each human culture “fills” it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us

feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through 're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns, nomenclature, and internal coherence of meaning and representation. In a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist

simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** The lower domain of experience is what we do not directly observe, the quantum realm, while the higher level is what we ordinarily observe, the classical realm. Based purely on experiment, the formulation of quantum theory initially placed the great divide between the quantum and classical domains between the measuring apparatus and the particle (Herbert, 1985). However, in the early 1930s, John von Neumann, in his rigorous mathematical treatise on quantum mechanics, found no support for such a division, compelling him to conclude that the wave for was collapsed by consciousness. Von Neumann's treatise has been called the "quantum bible" and "the most influential book on quantum theory ever written" (Herbert, 1985).

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t-T_9)F_9\big(R_9(t-T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t-T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) and the effective prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting $dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

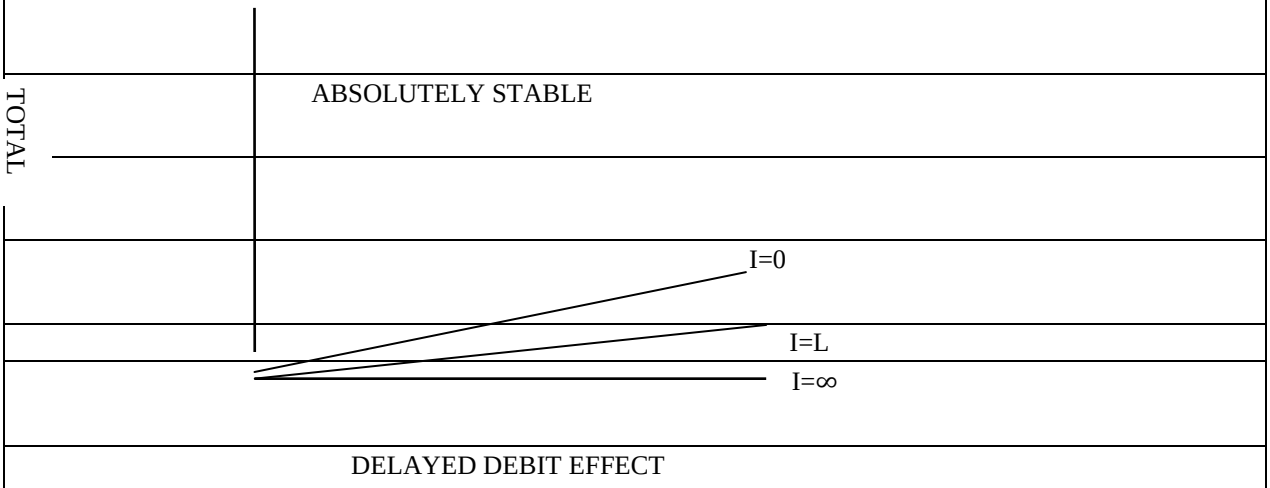



FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.

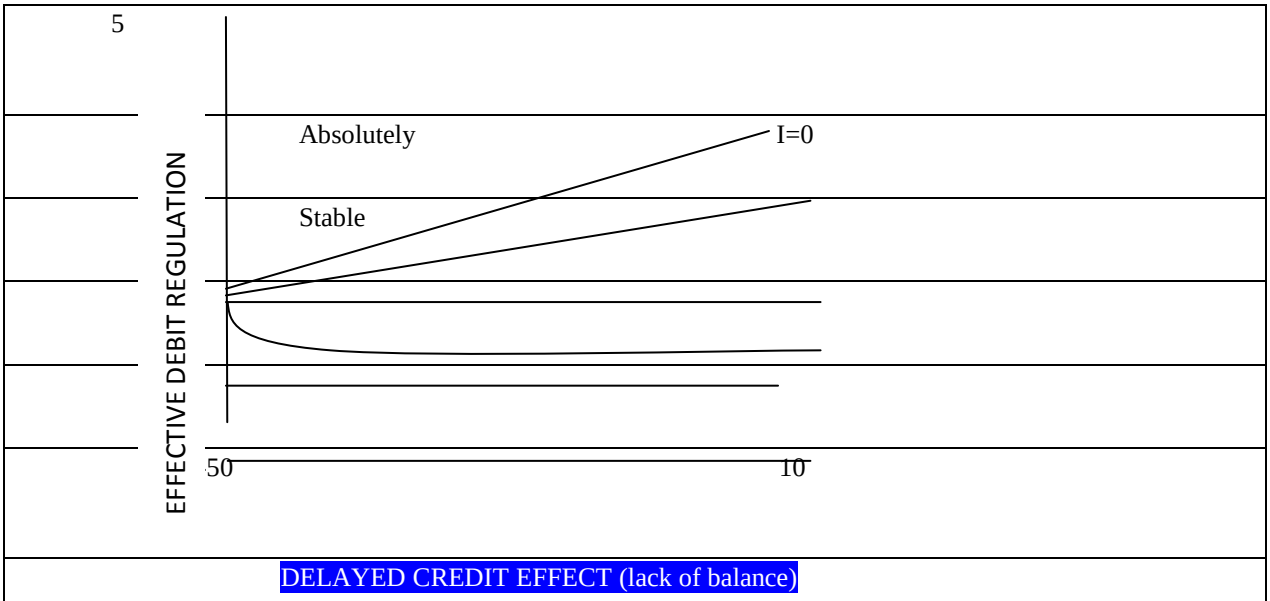



Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the

constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^\infty \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9\left(U_9 - V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9\right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$ then given the constraints of $f(T_9)$, $(A1.5)$ *can never be satisfied* Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDERATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 -$$

$$V_9 M_9 -]e - T_9\lambda_9 = 0$$

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9 Z_9 \sin(Z_9 T_9) - (A_9 M_9 + Y_9) + (A_9 M_9 + X_9)\cos(Z_9 T_9) = 0$$

$$(A_9 M_9) Z_9 - M_9 Z_9 \cos(Z_9 T_9) + (A_9 M_9 - X_9)\sin(Z_9 T_9) = 0$$

Where $A_9 = U_9 - V_9$ Using (A2.3), substitute for $Z_9^2$ and $Z_9$ in (A2.2):

$W_9^2 (X_9^2 - 3 M_9 A_9 X_9 - M_9^2 (X_9 - Y_9) + 2 M_9^3 A_9 + 2 A_9^2 M_9^2]$

$$+W\left[MA^{3}-A^{2}\mathrm{X}+2\mathrm{MA}(2\mathrm{X}+\mathrm{Y})-2X^{2}\right]+A^{2}(\mathrm{X}+\mathrm{Y})=0$$

*where* $W_9$=1-cos ($\mathrm{Z}_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (\mathrm{X}_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ \mathrm{X}_9 - Y_9/2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range 0⍰$W_9$ ⍰2 *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (\mathrm{X}_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\ \mathrm{M}_9)^2\ (2A_9\mathrm{M}_9 - (\mathrm{X}_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9M_9^2 - 3M_9X_9\right) + W_9\left(3M_9A_9^2 - 2A_9X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9 + Y_9\right)$$

$=2M_9^2W_9^2(M_9 + A_9) + M_9A_9^2W_9 + \left(2A_9M_9 - X_9 + Y_9\right)(A_9W_9 + M_9W_9^2$

$+ (2- W_9)((2M_9X_9 + M_9Y_9)W_9 + A_9(X_9 + Y_9))$,

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range 0 $W_9$ 2.Defining

$Q_9(W_9, A_9) = a_9W_9^2 + c_9$ *where* $a_9$ $b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

**Harald Atmanspacher, Hans Primas** show that such an epistemic perspective can be consistently complemented by Einstein's ontically oriented position. **Time, Quantum and Information 2003, pp 301-321 Epistemic and Ontic Quantum Realities Harald Atmanspacher, Hans Primas Robert W. Spekkens** present a toy theory that is based on a simple principle: the number of questions about the physical state of a system that are answered must always be equal to the number that are unanswered in a state of maximal knowledge. Many quantum phenomena are found to have analogues within this toy theory. These include the noncommutativity of measurements, interference, the multiplicity of convex decompositions of a mixed state, the impossibility of discriminating nonorthogonal states, the impossibility of a universal state

inverter, the distinction between bipartite and tripartite entanglement, the monogamy of pure entanglement, no cloning, no broadcasting, remote steering, teleportation, entanglement swapping, dense coding, mutually unbiased bases, and many others. The diversity and quality of these analogies is taken as evidence for the view that quantum states are states of **incomplete knowledge** rather than states of reality. A consideration of the phenomena that the toy theory **fails to reproduce, notably, violations of Bell inequalities** and the existence of a Kochen-Specker theorem, provides clues for how to proceed with this research program.DOI: http://dx.doi.org/10.1103/PhysRevA.75.032110 **Evidence for the epistemic view of quantum states A toy theory Phys Rev A 75, 032110 – Published 19 March 2007 Robert W. Spekkens In such specialized and autonomous areas, it is to the eminent and erudite authors to interpret the contradistinctions and contradictions' in the dove tailing explaining. This I feel would be yeoman service in the creation of Nature's general ledger. M. S. Leifer and O. J. E. Maroney** examine the relationship between quantum contextuality (in both the standard Kochen-Specker sense and in the generalized sense proposed by Spekkens) and models of quantum theory in which the quantum state is maximally epistemic. We find that preparation noncontextual models must be maximally epistemic, and these in turn must be Kochen-Specker noncontextual. This implies that the Kochen-Specker theorem is sufficient to establish both the impossibility of maximally epistemic models and the impossibility of preparation noncontextual models. The implication from preparation noncontextual to maximally epistemic then also yields a proof of Bell's theorem from an Einstein-Podolsky-Rosen-like argument. DOI: http://dx.doi.org/10.1103/PhysRevLett.110.120401 Received 25 August 2012Published 20 March 2013 © 2013 American Physical Society **Maximally Epistemic Interpretations of the Quantum State and Contextuality Phys. Rev. Lett 110, 120401 – Published 20 March 2013 M. S. Leifer and O. J. E. Maroney.** Epistemological, onitic and ontological aspects particularly of quantum mechanics have invited

profound thinkers to contribute their mite to the growing literature on the subject, which includes the contributions from the subterranean realm and ceratoid dualism of philosophy like that Chaos theory of Deleuze, Derrida , Kierkegaard and others. **Andrei Khrennikov** shows that the Dirac-von Neumann formalism for quantum mechanics can be obtained as an approximation of classical statistical field theory. This approximation is based on the Taylor expansion (up to terms of the second order) of classical physical variables – maps f: Ω → R, where Ω is the infinite-dimensional Hilbert space. The space of classical statistical states consists of Gaussian measures ρ on Ω having zero mean value and dispersion σ2 (ρ) ≈ h. This viewpoint to the conventional quantum formalism gives the possibility to create generalized quantum formalisms based on expansions of classical physical variables in the Taylor series up to terms of nth order and considering statistical states ρ having dispersion σ2(ρ) = h n (for n = 2 we obtain the conventional quantum formalism). **December 2005, Volume 18, Issue 7, pp 637-650 Date: 28 Nov 2005 Generalizations of Quantum Mechanics Induced by Classical Statistical Field Theory Andrei Khrennikov** It is widely accepted that consciousness or, in other words, mental activity is in some way correlated to the behavior of the brain or, in other words, material brain activity. Since quantum theory is the most fundamental theory of matter that is currently available, it is a legitimate question to ask whether quantum theory can help us to understand consciousness. Several approaches answering this question affirmatively, proposed in recent decades, will be surveyed. It will be pointed out that they make different epistemological assumptions, refer to different neurophysiological levels of description, and adopt quantum theory in different ways. For each of the approaches discussed, these imply both problematic and promising features which will be indicated. **Discrete Dynamics in Nature and Society Volume 2004 (2004), Issue 1, Pages 51-73 http://dx.doi.org/10.1155/S102602260440106X Quantum theory and consciousness: an overview with selected examples Harald Atmanspacher Copyright © 2004 Hindawi Publishing**

**Corporation. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.** Another thought provoking and investigatory article is by Peter G. Lewis et al. Perhaps the quantum state represents information about reality, and not reality directly. Wave function collapse is then possibly no more mysterious than a Bayesian update of a probability distribution given new data. **Peter G. Lewis et al** consider models for quantum systems with measurement outcomes determined by an underlying physical state of the system but where several quantum states are consistent with a single underlying state—i.e., probability distributions for distinct quantum states overlap. Significantly, they demonstrate by example that additional assumptions are always necessary to rule out such a model. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.150404 **Distinct Quantum States Can Be Compatible with a Single State of Reality Phys. Rev. Lett. 109, 150404 – Published 9 October 2012 Peter G. Lewis et al** Communication complexity of a quantum channel is the minimal amount of classical communication required for classically simulating a process of state preparation, transmission through the channel and subsequent measurement. It establishes and perpetuates a limit on the power of quantum communication in terms of classical resources. **Alberto Montina** shows that classical simulations employing a finite amount of communication can be derived from a special class of hidden variable theories where quantum states represent statistical knowledge about the classical state and not an element of reality. This special class has attracted strong interest very recently. The communication cost of each derived simulation is given by the mutual information between the quantum state and the classical state of the parent hidden variable theory. Finally, find that the communication complexity for single qubits is smaller than 1.28 bits. The previous known upper bound was 1.85 bits. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.110501 **Epistemic View of Quantum States and**

**Communication Complexity of Quantum Channels Phys. Rev. Lett 109, 110501 – Published 12 September 2012 Alberto Montina for models see elsewhere .** Stability analysis of what? Quantum state representing reality or our knowledge of reality? Does the quantum state represent reality or our knowledge of reality? In making this distinction precise, **Nicholas Harrigan, Robert W. Spekkens** are led to a novel classification of hidden variable models of quantum theory. They show that representatives of each class can be found among existing constructions for two-dimensional Hilbert spaces. Approach also provides a fruitful new perspective on arguments for the nonlocality and incompleteness of quantum theory. Specifically, **Nicholas Harrigan, Robert W. Spekkens** show that for models wherein the quantum state has the status of something real, the failure of locality can be established through an argument considerably more straightforward than Bell's theorem. The historical significance of this result becomes evident when one recognizes that the same reasoning is present in Einstein's preferred argument for incompleteness, which dates back to 1935. This fact suggests that Einstein was seeking not just any completion of quantum theory, but one wherein quantum states are solely representative of our knowledge. Hypothesis is supported by an analysis of Einstein's attempts to clarify his views on quantum theory and the circumstance of his otherwise puzzling abandonment of an even simpler argument for incompleteness from 1927. **February 2010, Volume 40, Issue 2, pp 125-157 Date: 09 Jan 2010 Einstein, Incompleteness, and the Epistemic View of Quantum States Nicholas Harrigan, Robert W. Spekkens** statistical model of the probabilistic description of physical reality has been studied by many authors. **Andrei Khrennikov posits** a contextual statistical model of the probabilistic description of physical reality. Here contexts (complexes of physical conditions) are considered as basic elements of reality. There is discussed the relation with QM. He proposes a realistic analogue of Bohr's principle of complementarity. In the opposite of the Bohr's principle, **Andrei Khrennikov** principle has no direct relation with mutual exclusivity for

observables. To distinguish his principle from the Bohr's principle and to give better characterization, he changes the terminology and speaks about supplementarity, instead of complementarity. Supplementarity is based on the interference of probabilities. It has quantitative expression trough a coefficient which can be easily calculated from experimental statistical data. There is need not appeal to the Hilbert space formalism and noncommutativity of operators representing observables. Moreover, in our model there exists a pair of supplementary observables which cannot be represented in the complex Hilbert space. There are discussed applications of the principle of supplementarity outside quantum physics. **October 2005, Volume 35, Issue 10, pp 1655-1693 The Principle of Supplementarity: A Contextual Probabilistic Viewpoint to Complementarity, the Interference of Probabilities and Incompatibility of Variables in Quantum Mechanics Andrei Khrennikov Foundations of Physics** Description of the vacuum in Yang-Mills theory and arrive at a physical interpretation of the pseudoparticle solution and the attendant violation of symmetries has been done by **R. Jackiw and C. Rebbi.** It is here vaccuum stability forms the bastion, pillar, post and stylobate of the mechanisms of yang mills Theory and action of a particle. The existence of topologically inequivalent classical gauge fields gives rise to a family of quantum mechanical vacua, parametrized by a CP-nonconserving angle. The requirement of **vacuum stability** against gauge transformations renders the vacua chirally noninvariant. DOI: http://dx.doi.org/10.1103/PhysRevLett.37.172 **Vacuum Periodicity in a Yang-Mills Quantum Theory Phys. Rev. Lett 37, 172 – Published 19 July 1976 R. Jackiw and C. Rebbi.** Three-dimensional Yang-Mills and gravity theories augmented by gauge-invariant mass terms are analyzed. These topologically nontrivial additions profoundly alter the particle content of the models and lead to (eb) quantization of a dimensionless mass-coupling-constant ratio. The vector field excitations become (=) massive, with (e&eb) spin 1 (rather than massless with spin 0), and the mass provides (eb) an infrared cutoff. The gravitation

acquires (e) mass, mediates (e&eb) finite-range interactions, and has (e) spin 2 (rather than being absent altogether); although its mass term is of (e)third derivative order, there are no ghosts or acausalities. DOI: http://dx.doi.org/10.1103/PhysRevLett.48.975 © 1982 The American Physical Society **Three-Dimensional Massive Gauge Theories Phys. Rev. Lett 48, 975 – Published 12 April 1982 S. Deser et al** All scattering amplitudes in (eb) the maximally supersymmetric N=4 super-Yang–Mills theory possess (e) a new, dual superconformal symmetry which extends the previously discovered dual conformal symmetry of MHV amplitudes. To reveal this property authors formulate the scattering amplitudes as (=) functions on (e&eb) the appropriate dual superspace. Rewritten in this form, all tree-level MHV and next-to-MHV amplitudes exhibit manifest dual superconformal symmetry. Authors propose a new, compact and Lorentz covariant formula for (e) the tree-level NMHV amplitudes for (e) arbitrary numbers and types of external particles. The dual superconformal symmetry is broken at loop level by infrared divergences. However, authors provide evidence that the dual conformal anomaly of (e) the MHV and NMHV superamplitudes is (=) the same and, therefore, their ratio is (=) dual conformally invariant. Authors show this explicitly for (e) the six-particle amplitudes at one loop. Authors conjecture that these properties hold for all, MHV and non-MHV, superamplitudes in N=4 SYM both at (eb) weak and at strong coupling. **Nuclear Physics B Volume 828, Issues 1–2, 21 March 2010, Pages 317–374 Dual superconformal symmetry of scattering amplitudes in N=4 super-Yang–Mills theory J.M. Drummonda DOI: 10.1016/j.nuclphysb.2009.11.022.** By comparison with numerical results in (eb) the maximal Abelian projection of lattice Yang–Mills theory, it is argued that the nonperturbative dynamics of (e) Yang Mills theory can be described by (e) a set of fields that take their values in (eb) the coset space SU(2)/U(1). The Yang–Mills connection is parameterized in (eb) a special way to separate (e&eb) the dependence on the coset field. The coset field is then regarded as (=) a collective variable and a method to obtain its effective

action is developed. It is argued that the physical excitations of (e) the effective action may be (=) knot solitons. A procedure to calculate the mass scale of knot solitons is discussed for (e) lattice gauge theories in the maximal Abelian projection. The approach is extended to (e&eb) the SU (N) Yang–Mills theory. A relation between the large N limit and the monopole dominance is pointed out. **Physics Letters B Volume 458, Issues 2–3, 8 July 1999, Pages 322–330 An effective action for monopoles and knot solitons in Yang–Mills theory Sergei V. Shabanov DOI: 10.1016/S0370-2693(99)00612-7.** It is here the quintessential formulation of consciousness occurs. Like a man standing on the threshold of infinity trying to ponder what lies beyond the veil, which separates the seen from unseen a true seeker tries to lift the individual general ledger to the cosmic general ledger. Like in Darwinian evolution, mind also has to evolve from individual consciousness to cosmic consciousness. In Kaishmiri Shaivism one withdraws from the finite to the infinite, but one also goes on an outward journey from the infinite to the finite, because both the finite and the infinite have an intimate connection. The finite is not seen as unreal, but as a symbol of the infinite. There is no real distinction between them. Those two movements constitute Spanda, a key concept in Kaishmiri Shaivism. Spanda is the pulsation of the Absolute in different phases of being. There are no opposites like subject and object, unity or duality, absolute or relative. They are just different phases of the universal vibration of the Absolute. The goal is to realize or be (=) at once infinite and finite. One does not turn away from appearances (like in Advaita Vedanta), but one realizes that the Absolute manifests all things. Spanda, the eternal pulsation of the Absolute, oscillates between a passion to create and dispassion from the created. Through it the Absolute transform itself into all things and then returns back into the emptiness of its undifferentiated nature. In Kaishmiri Shaivism the Absolute is seen as pure consciousness (=being). The Absolute is an eternal all-pervasive principle, the highest reality, the nature of all entities eternally and blissfully at rest within its own nature. The Absolute is the nature of the Self (and

thus of us all). The Absolute is divine; it is Shiva, the Lord of the Universe. It is full of (e) conscious activity through (e&eb)which it generates the universe, and reabsorbs it into (e&eb) itself at the end of each cycle of creation. Thus we speak of monism, as everything resides within (eb) this one absolute consciousness. It sustains (eb) all things, it embraces (e) all things, and it pervades (e&eb) all things. All things are appearances within (eb) the absolute consciousness, but nevertheless real (in contrast to (e) Advaita Vedanta where (e) appearances are seen as unreal or illusionary). All things appear external (out there, outside ourselves), but they do not have (e) a being on their own. They do not exist as (=) separate entities on their own. Everything is contained within (eb) consciousness. What we see as objects are (=) manifestations of consciousness. The event which constitutes (eb) the universe is always internal events happening within (e&eb) consciousness because (e) their essential nature is consciousness itself. If a physical object were totally material, and independent or external to consciousness, it could never be experienced. The universe and (e&eb) consciousness are two aspects of a whole. The universe is (=) an attribute of consciousness which bears consciousness as (=) its substance. Doctrine of Vibration: S.G.Dyczkowski what you "see" is (=) not always what you "get." Many people mistakenly take their own visions literally without (e) "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in (eb) the physical sense but in a psychological sense. The archetypes of (e) the collective unconscious are arrayed behind (e) the scenes of current worldwide conditions, of (e) crisis and confusion. They mirror (e&eb) our own states back at us, whether we perceive them as (=) such or interpret them plausibly or not (Note the essay has been written in model able form and models are given in various papers.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4\,a_9 \square - b_9 \square 0$$

And

$$b_9^{\,2} \square 4\,a_9 c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9\, in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\, X_9$ ,

$Let\ Y_9 = \ y_9\, X_9$And try the substitution $A_9 = X_9/(2 + A_9\ y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + \ n_9 y_9)$ we have

$$W_9^{\,2}\, y_9 \left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^{\,2}\left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9\, X_9 \left[\frac{X_9 (1 + n_9 y_9)}{M_9^{\,2}(2 + n_9 y_9)^2} + 2\, y_9 \left(n_9 - 1\right) + \frac{X_9^{\,2}(9 + y_9)}{M_9^{\,2}(2 + n_9 y_9)} = 0.\ (A2.12)\right.$$

Since $n_9$ and $A_9$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9 = (2 + n_9 y_9)$, *the condition* (A2.99?) *can be written as*

$[\frac{X_9(E_9-9)}{E_9^2}]^2[\frac{X_9^2}{M_9^2} - 4\, y_9\ E_9^2 \frac{X_9}{M_9^2} + 4y_9 (n_9 - 2)E_9^3]$ .

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9 / W_9^2)$ can be expressed as

$2\, E_9 [y_9 E_9 + (y_9\ E_9 (9 + y_9)(E_9 - n_9))^{9/2}] > \frac{X_9}{M_9^2} > 2E_9 [y_9 E_9 - y_9 E_9 (9 + y_9)(E_9 - n_9))^{9/2}]$

Concept Of Being From Shaivite Point Of View: In His Famous Work Entitled On The Conservation Of Force (1847), Hermann Von Helmholtz Had Divided The Totality Of The Energy In The Universe Between Potential Energy And Kinetic Energy And Affirmed The Reciprocal Convertibility Of The Two. In 1852, William Thomson Pointed Out That There Exists A Sub-Ensemble Within Kinetic Energy, Heat, Which, Once It Has Been Generated, Is No Longer Entirely Convertible Into Potential Energy-Or Into Any Other Form Of Kinetic Energy. Considering That The (Partial) Reconversion Of Heat Into Labor Is Possible Only In Situations That Present A Disparity In Temperature, And That Heat Tends To Pass From Warmer To Cooler Bodies By Spreading On An Even Temperature Level Through Space, Thomson Concluded That The Universe Tends Towards A Final State Where Any Energetic Transformations, Every Movement And Every Form Of Life Will Cease: The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO: The Ultimate Nature Of Anything Cannot Change. Change Of Any Sort Is Merely Apparent (Vivarta). The World Of Change And Becoming Is A False Super- Imposition (Adhyaropa, Adhyasa) On The Absolute. It Appears So On The Screen Of Individual, Collective Consciousness, Albeit One Cannot Be So Sure Of Cosmic, For There Are Many Reasons That Cosmic Transactions Also Had To Be Subjected To Change

By ParaBrahaman. In Cosmic Terms, The Mistake (Bhranti) Consists Of The Supposition That The Real Brahman Is The Unreal Universe And The Unreal Universe Is The Real Brahman. In Microcosmic Terms, It Is The Mistake Of Falsely Conceiving The Body, Mind Or Even One's Personality To Be The Self. Note We Have Taken Self As Witness Consciousness. Notwithstanding There Is Nothing Wrong In The Conception Of Self As A Dynamical Presupposition, In That Individual Consciousness And Collective Consciousness And Even Cosmic Consciousness Evolutes In To ParaBrahaman. In The Same Way As The Image Of A Snake Is Falsely Superimposed On A Rope, Similarly The Universe Is Falsely Projected Onto The Real Substratum, The Brahman. Ignorance Is Not Merely A Personal Lack Of Knowledge, But A Cosmic Principle. As Such It Is Called "Maya," The Indefinable Factor (Anirvacaniya) That Brings This Mistake In Identity About. The Reality Status Of This Cosmic Illusion Is Also Indefinable: On The One Hand It Is Not Brahman, The Sole Reality; On The Other Hand It Is Not Absolutely Non-Existent Like A Hare's Horn Or The Son Of A Barren Woman. Note Here The Vedantin Avoids Talking About Zero Which We Have Avidly And Assiduously Stressed About. Nothing Is The Factor That Exists. Also See the Remarks about Satre's Being And Nothingness. Brahman Is The Source Of World Appearances Only In The Sense Of Being Their Unconditioned Ground Or Essential Nature. The Universe Is False Not Because It Has No Nature Of Its Own But Because It Does Have One. The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski. Concept Of Being From Shaivite Point Of View: It Is The 'I Am The Body/Mind' Belief That Gives Rise To The 'I Am Not The World' Belief. These Two Beliefs Are (=) Co-Created."Consciousness Projects The Appearance Of The Mind, Body And World By Taking (Eb) The Shape Of Thinking, Sensing And Perceiving."“Attention Is Consciousness With An Object. When The Object Vanishes, Attention Simply Remains What It Always Is Consciousness."“There Is No Purpose To Meditation. The Purpose Is Already Accomplished."“Everything

That Is Experienced Is Experienced By, Through, In And As Consciousness.”“The Seen Cannot Be Separated From Seeing And Seeing Cannot Be Separated From Consciousness.” “**The Reality Of Any Experience Is Not Hidden In The Appearance; It Is Expressed By The Appearance The Transparency Of Things**' By Rupert Spira Book Review By Dennis Waite The Following Is A Review Of Rupert's Book: 'The Transparency Of Things: Contemplating The Nature Of Experience’. There Is An Essay From The Book Here And The Essay On Cézanne May Be Read At The Awakened Eye Website Nature Rejects The Naiveté That Seeks Absolute Truth. We Are Beginning To Realize, Individually And Culturally, That "Realities" Are All Human Constructions. The Task Becomes One Of "Catching Ourselves In The Act" Of Creating Our Own "Reality" From The Flow Of Events. Human Truth Is Always An Engagement Of Mind With Experience. The Challenge Of The Therapist In These Times Of Chaotic Change Is To Validate The Concept That We Don't Have To Fear The Collapse Of What We Think We Are. Strange Attractors: Transference, Holography, And An **Archetype Burke, J. (2003). Strange Attractors: Transference, Holography, And An Archetype (Doctoral Dissertation, Pacifica Graduate Institute, 2003).** When we get to the topic and understanding of the Transcendental Logic, the brilliant and baffling Transcendental Deduction (TD) Recall the two movements just discussed, the one from experience to its conditions and the one from the forms of valid inference to the concepts that we must use in all judging (the Categories). This duality led Kant to his famous question of right (quid juris) (A84=B116): with what right do we apply the Categories, which are not acquired from experience, to the contents of experience? (A85=B117). Kant's problem here is not as arcane as it might seem. It reflects an important question: How is it that the world as we experience it conforms to our logic? In briefest form, Kant thought that the trick to showing how it is possible for the Categories to apply to experience is to show that it is necessary that they apply (A97).TD has two sides, though Kant never treats them separately John Marlado,

accounts, attributes, and ascribes the philosophy of Nishida KitarÅ is lucid, insightful, and deeply informative despite the highly questionable argument that gives the book its structure. In comparison with other, technically more accurate accounts of Nishida's philosophy, it is broader in scope than most but more single-minded in its approach. The analysis is for the most part limited to four books that have been translated into English, out of approximately 16 volumes of philosophical essays in Japanese, but it aims to capture the contours and most important details of Nishida's entire philosophy. To appreciate the achievement of this account, one would need only to read, in English translation or in the Japanese original, any of the volumes with their streams of repetitive, puzzling expressions and circuitous, often dead- end argumentation. Wilkinson's elucidation of details in the passages he samples makes it easier to understand why Nishida is considered Japan's greatest academic philosopher. A critical look at the frame of the book and some of its own claims reveals where, to my mind, it has gone astray. The argument that frames this book is repeated throughout: Nishida's philosophy is an attempt to articulate, in a characteristically Western philosophical manner, the Zen experience he had as a young man. Nishida first made the attempt by using the conceptual frameworks of Western philosophers, and after they failed him he developed his own, largely by recasting Buddhist insights. In the end, a philosophical articulation of the Zen worldview in Western terms or even in Eastern terms recast by conceptual, Western philosophical methods approved impossible. What Nishida attempted is incommensurable with Western philosophy. While the author first offers these statements as hypotheses (p. 2), one premise of his argument that there is such a thing as "the Zen world-view" or "the Zen experience “remains unquestioned; and another that Nishida's pivotal experience was the Zen experienceâis asserted as "beyond question" (p. 151; see also pp. 28, 48, 10, 114).Wilkinson's presentation of Zen in the first chapter is drawn largely from D. T. Suzuki, with help from Suzuki's portage Abe Masao; some comments on time, birth, and death in; and a few other sources. The

result is a caricature that imagines Zen as an invariable, unitary experience. In fact, Dagan's teachings on practice as the manifestation of enlightenment do not sit well with D. T. Suzuki's emphasis on a satori experience, and Suzuki does not enlighten us about what Zen in its various practices has historically been. The author would have benefited from acquaintance with other presentations by Zen teachers, such as that of ShunryÅ Suzuki'sZen Mind Beginner's Mind, with historical accounts of Zen traditions, and with contrasts between DÅgen's view of time and Sam Â· sÄraand that of other Buddhist philosophers. A little research in contemporary Zen scholarship would have shown how implausible is the talk of "nirvana or awareness of mu" (p. 24), "the basic Zen position [that] never varies" (p. 3), "the Zen vision" (pp. 5, 7), "Zen epistemology" , "Zen experience" (p. 59), "the conception of time involved in Zen" (p. 15), and "the fundamental Zen assertion that there is no absolute time" (p. 16). Even if DÅgen's notion that time and being are nondual (p. 16) epitomized most of Zen, Nishida's early notion of "a transcendent, unchanging reality apart from time" (p. 51) would contrast sharply with it. Equally questionable is Wilkinson's contention that Nishida's "pure experience" is equivalent to Zen satori. The notion of pure experience Nishida developed in his first major work may be pivotal for his entire philosophy, but he never claims or implies that any part of his philosophy presupposes an exceptional experience like satori that grounds his convictions and renders them unverifiable to those who have not had the experience, as Wilkinson intimates (pp. 55, 159, and 161). Nishida presents his notions of pure experience and intellectual intuition, the grasp of its unity, and later... Perennialism and Constructivism is studied by Randolph T Dible : experiencer or the observer and the transcendental subjectivity itself Pure objectivity is essentially a deconstruction, apophatic and negativa The encountering of the experiencer or observer—transcendental subjectivity itself—at the foundation of the world leads inevitably to the recognition of pure objectivity as ultimate reality (which can be taken as its ultimate deconstruction, analogous to the apophatic or via

negativa), from which objects derive their value, weight, significance, meaning or objectivity. In this way, pure objectivity can be seen as the supra-self-evident Axiological Axiom, so to speak, even Unconditional Love, in romantic terms. This axiology (value theory) has a structure inverse to the relationship between transcendental subjectivity as the radical unity of pure self-reference and on the other hand, the world of forms, as mere traces (representations, indications) of the unique, original "first distinction" Spencer-Brown speaks of at the foundation of his calculus. That is, all forms (i.e., distinctions, differences) would reduce to being the first distinction, also known as the marked state, which can be called penultimate reality (pure self-reference or transcendental subjectivity: the Spirit which animates us), except that forms are complimentary to their content, which is their objectivity or value, which would reduce to the unmarked state or ultimate reality. It is the incongruity of form (thoughts; Whitehead's "negative prehensions") and value (feelings; Whitehead's "positive prehensions," or my notion of objectivity, meaning and qualia; in short, the non-formal aspects of experience) that holds forms open and keeps them from absolute reduction. This accounts for the brute, concrete persistence of the "functional illusion"-- to use a term from Dzogchen Buddhism-- of the world. Thus this system has an axiology of metaphysical objectivity grounded on the ideal of pure objectivity as the source of all value, meaning and significance, itself the very fecundity of profundity, which is the motive of drawing the distinction in the first place **(Wikipedia, Kant's writings and Logic Of Sense)** Nano rings were created by accident while intending to make quantum dots. They have interesting optical properties associated with excitons and the Aharonov–Bohm effect. Application of these rings used as light capacitors or a buffer includes photonic computing and communications technology. Analysis and measurement of geometric phases in mesoscopic rings is ongoing. Several experiments, including some reported in 2012, show Aharonov-Bohm oscillations in charge density wave (CDW) current versus magnetic flux, of dominant period h/2e through CDW rings up to 85 **μm** in

circumference above 77 K. This behavior is similar to that of the superconducting quantum interference devices (seeSQUID). **Aharonov–Bohm nano rings (Wikipedia) Stability analysis of nano rings would help progressively accentuate the research about the materials division that uses nano rings.** From the profound contemplation and wisdom of the Buddha and Mahabodhisattvas, we know there is no such reality. Instead, egolessness (non-self) is the only path to understand the reality of the deluded life. All existences are subject to the law of causes and conditions. These include the smallest particles, the relationship between the particles, the planets, and the relationship between them, up to and including the whole universe! From the smallest particles to the biggest matter, there exists no absolute independent identity. Egolessness (non-self) implies the void characteristics of all existence. Egolessness (non-self) signifies (eb) the non-existence of permanent identity for self and existence (dharma). Sunyata stresses (eb) the voidness characteristic of (e) self and existence (dharma). Sunyata and Egolessness possess (e) similar attributes. As we have discussed before, we can observe the profound significance of Sunyata from (e) the perspective of inter-dependent relationships. Considering dharma-nature and (e&eb) the condition of nirvana, all existences are immaterial and of (e) a void-nature. Then we see each existence as (=) independent of each other. But then we cannot find any material that does exist independent of (e) everything else. So Egolessness also implies (eb) void-nature! From the observation of all existences, we can infer (eb) the theory of nirvana and the complete cessation of (e) all phenomena. From the viewpoint of phenomena, all existences are so different from (e) each other, that they may contradict (e&eb) each other. They are (=) so chaotic. In reality, their existence is (=) illusionary and arises from (e) conditional causation. They seem to exist on one hand, and yet do not (e) exist on the other. They seem to be united, but yet they are so different to (e) one another. They seem to exist and yet they do cease! Ultimately everything will return to (e) harmony and complete calmness. This is the nature of all existence. It is the

final resting place for all. If we can understand this reality and remove our illusions, we can find this state of harmony and complete calmness. **Teachings in Chinese Buddhism (6) sunyata (emptiness) in the Mahayana context (Wikipedia) Many models have been given concatenating Buddhist doctrines and Vaccum energy. These models essentially throw light on the essence of self(witness consciousness) and the cosmic general ledger (cosmic consciousness) while the imbalance in the gratification deprivation complex always remain unbalanced when it comes to individual general ledger , advocating the evolutionary theory of individual consciousness to cosmic consciousness. It is to be noted that it is not the question of knowing all the transactions but the wisdom of what lies behind the debit- Credit, Debit-Debit, or Credit-Credit principle that is most important for assimilation.** Look at this simple thesis to prove the point stated in the foregoing. Thesis: 'The World Has A Beginning In Time, And In Space It Is Also Enclosed In Boundaries.' Proof: 'For If One Assumes That The World Has No Beginning In Time, Then Up To Every Given Moment In Time An Eternity Is Elapsed, And Hence An Infinite Series Of States Of Things In The World, Each Following Another, Has Passed Away. But Now The Infinity Of A Series Consists Precisely In The Fact That It Can Never Be Completed Through A Successive Synthesis. Therefore An Infinitely Elapsed World-Series Is Impossible, So A Beginning Of The World Is A Necessary Condition Of Its Existence, Which Was The First Point To Be Proved (Kant, Critique Of Pure Reason, Tr. By R. Guyer And A. W. Wood (Cambridge: Cambridge University Press, 1998): B 454, P. 470). As Quoted In The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO"The Cyclical Hypothesis, So Heavily Criticized By Nietzsche (VP II 325 And 334), Arises In This Way." In Fact, Nietzsche Was Not Criticizing The Cyclical Hypothesis But Only The Particular Form Of That Hypothesis Presented In Vogt's Work. All Of Nietzsche's Texts Without Exception Speak Of The Eternal Return As The Repetitio Of The Same Events Within A Cycle Which Repeats Itself

Eternally. If Deleuze's Interpretation Holds That The Eternal Return Is Not A Circle, Then What Is It? A Wheel Moving Centrifugally, Operating A "Creative Selection," "Nietzsche's Secret Is That The Eternal Return Is Selective" Says Deleuze: The Eternal Return Produces Becoming-Active. It Is Sufficient To Relate The Will To Nothingness To The Eternal Return In Order To Realize That Reactive Forces Do Not Return. However Far They Go, However Deep The Becoming-Reactive Of Forces, Reactive Forces Will Not Return. The Small, Petty, Reactive Man Will Not Return. Affirmation Alone Returns, This That Can Be Affirmed Alone Returns, Joy Alone Returns. Everything That Can Be Denied, Everything That Is Negation, Is Expelled Due To The Very Movement Of The Eternal Return. We Were Entitled To Dread That The Combinations Of Nihilism And Reactivity Would Eternally Return Too. The Eternal Return Must Be Compared To A Wheel; Yet, The Movement Of The Wheel Is Endowed With Centrifugal Powers That Drive Away The Entire Negative. Because Being Imposes Itself On Becoming, It Expels From Itself Everything That Contradicts Affirmation, All Forms Of Nihilism And Reactivity: Bad Conscience, Ressentiment..., We Shall Witness Them Only Once. [...] The Eternal Return Is The Repetition, But The Repetition That Selects, The Repetition That Saves. Here Is The Marvelous Secret Of A Selective And Liberating Repetition. There Is No Need To Remind The Reader That Neither The Image Of A Centrifugal Movement Nor The Concept Of A Negativity-Rejecting Repetition Appears Anywhere In Nietzsche's Writings, And Indeed Deleuze Does Not Refer To Any Text In Support Of This Interpretation. Further, One Could Highlight That Nietzsche Never Formulates The Opposition Between Active And Reactive Forces, Which Constitutes The Broader Framework Of Deleuze's Interpretation. For Some Years, Marco Brusotti Has Called Attention To The Fact That Deleuze Introduced A Dualism That Does Not Exist In Nietzsche's Writings. To Be Sure, The German Philosopher Describes A Certain Number Of "Reactive" Phenomena (For Example, In The Second Essay Of The Genealogy Of Morality, § 11, He Talks About

"Reactive Affects" [Reaktive Affekte], "Reactive Feelings" [Reaktive Gefühlen], Reactive Men [Reaktive Menschen]); But These Are Nonetheless The Result Of Complex Ensembles Of Configurations Of Centers Of Forces That Remain In Themselves Active. Neither The Word Nor The Concept Of "Reactive Forces" Ever Appears In Nietzsche's Philosophy. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO Some Parts Are Deleted Due To Spatial Constraints. In the midst of thoughts about the eternal return we find at least two other thematic axes On the one hand, the view of the **world as a constant flux of forces without any goal, law, or rules of becoming. A chaos sive natura de-divinized and de-anthropomorphized which constitutes the "ontological substratum**" of the whole of Nietzsche's reflexions On the other hand, an ensemble of fragments of an anthropologico-sociological character, designing a path of liberation leading to the creation of superior individuals by way of a profound transformation of their instinctual structure. This transformation must be achieved by a practice of solitude and internal struggle towards the liberation from the ancient representations of the world and from the incorporated herd values. For an analysis of these thematic perspectives, see Paolo D'Iorio, La linea e il circolo. Cosmologia e filosofia dell'eterno ritorno in Nietzsche (Genova: Pantograf, 1995): 233-322. With most of the parameters quantified, with assumptions there must be picturesque view of the self and non self, space and time, science and philosophy with the concatenational consummations. After The Chapter "On Redemption," Where Zarathustra Dares Not Expose His Doctrine, The Eternal Return Begins To Be Enunciated In Part Three Of The Work. In The First Place, It Is The Dwarf Who Formulates It In The Chapter "On The Vision And The Riddle." Facing The "Gate Of The Instant" Which Symbolizes The Two Infinities That Stretch Towards The Past And The Future, The Dwarf Whispers: "All Truth Is Crooked, Time Itself Is A Circle" The Dwarf Represents The Spirit Of Gravity, And He Embodies The Herd Morality, "The Belittling Virtue" Which Is The Title Of Another Chapter From Part III. The Dwarf Can

Endure The Eternal Return Without Great Difficulties Because He Has No Aspirations; Unlike Zarathustra He Does Not Wish To Climb The Mountains That Symbolize Elevation And Solitude. In Two Unpublished Notes, From The Summer And The Fall Of 1883, Nietzsche Writes: The Doctrine Is At First Favored By The Rabble, Before It Gets To The Superior Men. The Doctrine Of Recurrence Will First Smile To The Rabble, Which Is Cold And Without Any Strong Internal Need. It Is The Most Ordinary Of Life Instincts, Which Gives Its Agreement First. Hence, The Content Of The Doctrine Is The Same, But Whereas The Dwarf Can Endure It (Because He Interprets It According To The Pessimistic Tradition For Which "Nothing Is New Under The Sun"), Zarathustra, Who Is The "Advocate Of Life" Regards The Eternal Return As The Strongest Objection To Existence, And As The Rest Of The Dream Suggests, He Does Not Yet Succeed In Accepting It After The Vision At The Gate Of The Instant, The Chapter Is Brought To An End By The Enigma Of The Shepherd. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO .How does the memory affect the GD complex or for that matter the gratification deprivation of an individual or a collective identity? Non-Markovian local-in-time master equations give a relatively simple way to describe the dynamics of open quantum systems with memory effects. Despite their simple form, there are still many misunderstandings related to the physical applicability and interpretation of these equations. Here, we clarify these issues both in the cases of quantum and classical master equations. Further introduction of the concept of a classical non-Markov chain signified through negative jump rates in the chain configuration. **E-M Laine et al 2012 J. Phys. B: At. Mol. Opt. Phys. 45 154004 doi:10.1088/0953-4075/45/15/154004 Local-in-time master equations with memory effects: applicability and interpretation.** Supposing That There Were Indeed An "Energy Of Contraction" Constant In All Centers Of Force Of The Universe, It Remains To Be Explained Where Any Difference Would Ever Originate. It Would Be Necessary For The Whole To Dissolve Into An Infinite Number Of

Perfectly Identical Existential Rings And Spheres, And We Would Therefore Behold Innumerable And Perfectly Identical Worlds COEXISTING [Nietzsche Underlines This Word Twice] Alongside Each Other. Is It Necessary For Me To Admit This? Is It Necessary To Posit An Eternal Coexistence On Top Of The Eternal Succession Of Identical Worlds? The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. Science is usually portrayed as dehumanizing. Brave New World epitomizes this fear. "The more we understand the world, the more it seems completely pointless" (Steven Weinberg). Certainly science can seem chilling when conceived in the abstract as a metaphysical world-picture. We may seem to find ourselves living in a universe with all the human meaning stripped out: Participants in a soulless dance of molecules, or harmonics of pointlessly waggling superstrings and their braneworld cousins. Nature seems loveless and indifferent to our lives. What rights have we to be happy? **(Brave New World: Aldous Huxley).** Religion can afford to be dogmatic but science cannot. Look at this passage which counters both scientists and philosophers alike: However, Hartmann Objects That The Regressive Movement Postulated By Schopenhauer Is Possible Only In Thought: It Remains Nothing More Than An "Ideal Postulate" With No Real Object And Which "Does Not Teach Us Anything About The Real Process Of The World That Unfolds In A Movement Contrary To This Backwards Movement Of Thought" (Hartmann, Philosophie Des Unbewussten, Third Edition (1871): 772). Hartmann Affirms That If Unlike Schopenhauer One Admits The Reality Of Time And Of The World Process, One Must Also Admit That The Process Must Be Limited In The Past And Therefore That There Must Be An Absolute Beginning. In Hartmann's Mind, Failure To Do So Would Result In Positing The Contradictory Concept Of An Accomplished Infinity: "The Infinity That From The Point Of View Of Regressive Thinking, Remains An Ideal Postulate, Which No Reality May Correspond To, Must, For The World, Whose Process Is, On The Contrary, A Progressive Movement, Open Up To A Determinate Result; And Here The Contradiction Comes To Light" (Hartmann

(1871): 772). What Really "Comes To Light" In This Passage Is The Fact That Hartmann Does Not Provide A Demonstration But A Petitio Principii. Indeed, The Concept Of The World Process Analytically Contains The Concept Of A Beginning Of The World. In All Rigors, It Is Therefore Impossible To Demonstrate These Concepts With Reference To Each Other. Secondly, Hartmann's View That One Is Bound To Accept The Reality Of The World Process Even If One Rejects The Ideality Of Schopenhauer's Time Is Mistaken. Hartmann Believes That If Time Is Real There Must Be A World Process With Both An Absolute Beginning And An Absolute End. Without Any Justification, Hartmann Jumps From Schopenhauer's Negated Time To Oriented Time. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. In the last years several theoretical papers discussed if time can be an emergent property deriving from ) quantum correlations. Here, to provide an insight into how this phenomenon can occur, authors present an experiment that illustrates Page and Wootters' mechanism of "static" time, and Gambini et al. subsequent refinements. A static, entangled state between a clock system and the rest of the universe is perceived as evolving by internal observers that test the correlations between the two subsystems. They implement this mechanism using an entangled state of the polarization of two photons, one of which is used as a clock to gauge the evolution of the second: an "internal" observer that becomes correlated with the clock photon sees the other system evolve, while an "external" observer that only observes global properties of the two photons can prove it is static.Phys. Rev. A 89, 052122 (2014)DOI:10.1103/PhysRevA.89.052122 arXiv: 1310.4691 [quant-ph] **Time from quantum entanglement: an experimental illustration Ekaterina Moreva.** Time is an entanglement .When the new ideas of quantum mechanics spread through science like wildfire in the first half of the 20th century, one of the first things physicists did was to apply them to gravity and general relativity. The results were not pretty. It immediately became clear that these two foundations of modern physics were entirely

incompatible. When physicists attempted to meld the approaches, the resulting equations were bedeviled with infinities making it impossible to make sense of the results. Then in the mid-1960s, there was a breakthrough. The physicists John Wheeler and Bryce DeWitt successfully combined the previously incompatible ideas in a key result that has since become known as the Wheeler-DeWitt equation. This is important because it avoids the troublesome infinites—a huge advance. But it didn't take physicists long to realise that while the Wheeler-DeWitt equation solved one significant problem, it introduced another. The new problem was that time played no role in this equation. In effect, it says that nothing ever happens in the universe, a prediction that is clearly at odds with the observational evidence. This conundrum, which physicists call 'the problem of time', has proved to be thorn in flesh of modern physicists, who have tried to ignore it but with little success. Then in 1983, the theorists Don Page and William Wootters came up with a novel solution based on the quantum phenomenon of entanglement. This is the exotic property in which two quantum particles share the same existence, even though they are physically separated. Entanglement is a deep and powerful link and Page and Wootters showed how it can be used (e) to measure time. Their idea was that the way pair of entangled particles evolves is a kind of clock that can be used to measure change. But the results depend on how the observation is made. One way to do this is to compare the change in the entangled particles with an external clock that is entirely independent of the universe. This is equivalent to god-like observer outside the universe measuring the evolution of the particles using an external clock. In this case, Page and Wootters showed that the particles would appear entirely unchanging—that time would not exist in this scenario. But there is another way to do it that gives a different result. This is for an observer inside the universe to compare the evolution of the particles with the rest of the universe. In this case, the internal observer would see a change and this difference in the evolution of entangled particles compared with everything else is an important a measure of time. This is

an elegant and powerful idea. It suggests that time is an emergent phenomenon that comes about because of the nature of entanglement. And it exists only for observers inside the universe. **Time is an entanglement** (Wikipedia). Danish religious philosopher Soren Kierkegaard carried out a systematic critique of the pretensions of reason and an abstract rationalism which he believed that the modern age was nurturing. Condemning reflection as a "danger" that ensnares people in logical delays and machinations, Kierkegaard compared it to a prison. Reflection is for him a form of captivity, a bondage which "can only be broken by [passionate] religious inwardness" (1978: 81). **Reflection seduces individuals into thinking its possibilities are "much more magnificent than a paltry decision"** (1978: 82). It leads them to act **"on principle,"** to dwell on the deliberation of the context of their actions and the calculation of their worth or outcome. Kierkegaard argues that this drives away feeling, inspiration, and spontaneity, all of which are crucial for true inner being and a vital relation to God. For Kierkegaard, as Nietzsche would later agree, genuine inner being (and culture) is characterized by the tautness and tension of the soul which characterizes passionate existence. But the "coiled springs of life relationships ... lose their resilience" in reflection (1978: 78) and "everything becomes meaningless externality, devoid of [internal] character" (1978: 62). Kierkegaard thus contributes to the development of an irrationalist tradition that has echoes in some later postmodern thought. Kierkegaard might well have agreed with his contemporary Fyodor Dostoyevsky, who wrote: "An intelligent [reflective] man cannot seriously become anything ... excessive consciousness is a disease" (1974: 3, 5). In an age overtaken by rules and regulations, genuine action -- which Kierkegaard assumes to be subjective and spontaneous -- is frustrated at every turn. Complaining that we are too "sober and serious" (1978: 71) even at banquets, Kierkegaard bemoans the fact that even suicides are premeditated (1978: 68)! "That a person stands or falls on his actions is becoming obsolete; instead, everybody sits around and does a brilliant job of bungling through with the aid of some reflection

and also by declaring that they all know very well what has to be done" (1978: 73). Thus, it is passion, not reflection, that guarantees "a decent modesty between man and man [and] prevents crude aggressiveness" (1978: 62). "Take away the passion and the propriety also disappears" (1978: 64). The ambiguity in the word "passion" may cause some confusion here. To say that the age and its individuals are "passionless" is not to say there are no emotions whatsoever, but rather that there is no true spiritual inwardness and depth, no intensively motivated action and commitment. It suggests that passion exists only in a simulated, pseudo-form, "the rebirth of passion" through "talkativeness" (1978: 64). "Chattering" for Kierkegaard gets in the way of "essential speaking" and merely "reflects" inconsequential events (1978: 89-99). Hence, in "the present age," emotions -- which in fact are all too pronounced -- have been transformed into negative forces. Anticipating Nietzsche's genealogy of the "slave revolt" in morality, Kierkegaard claims that the "enthusiasm" of the prior age of Revolution, a "positively unifying principle," has become a vicious "envy," a "negatively unifying principle" (1978: 81), a leveling force in its own right insofar as those lacking in talent and resources want to tear down those who have them. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS By Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html.** In November 2011, the United States government argued before the US Supreme Court that it wants to continue utilizing GPS tracking of individuals without first seeking a warrant. In response, Justice Stephen Breyer questioned what this means for a democratic society by referencing Nineteen Eighty-Four. Justice Breyer asked, "If you win this case, then there is nothing to prevent the police or the government from monitoring 24 hours a day the public movement of every citizen of the United States. So if you win, you suddenly produce what sounds like 1984... "[65]. What is to be said in unmistakable terms is that when it comes to national security and innocent citizens there is question of compromise even it be an infringement and encroachment of so called

privacy which has been made a word to subterfuge the activities that are far said to be fair. Quantum theory has provoked intense discussions about its interpretation since its pioneer days. One of the few scientists who have been continuously engaged in this development from both physical and philosophical perspectives is Carl Friedrich von Weizsäcker. The questions he posed were and are inspiring for many, including the authors of this contribution. Weizsäcker developed Bohr's view of quantum theory as a theory of knowledge

## ABSOLUTE STABILITY

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range (0 W 9 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9 - 9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2 - 2y_9\frac{(E_9 + n_9^2 y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right) - 4y_9(E_9 - n_9)$$ 0

And has the solution

$$\frac{X_9}{M_9^2} > \frac{E_9}{(E_9 - 9)}[y_9(E_9 + n_9^2 y_9) + [y_9^2 (E_9 + n_9^2 y_9)^2 + 4y_9(E_9 - 9)(E_9 - n_9)]^{9/2}]$$

Given $E_9 > n_9$ this lower limit to $X_9 / M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9 / M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$ in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9 < n_9$ $(ie., 2A_9 M_9^2 < X_9 - Y_9)$ then (A2.9)and (A2.8)show that $Q_9$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2} > 2E_9[y_9 E_9 - (y_9 E_9 (9 + y_9)(E_9 - n_9)]^{9/2}]$$

And

$$E_9 > n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2 /2M_9 (X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$ and $X_9$ and $Y_9$ by $W_9^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9$ *and* separate (A3.1)into real and imaginary parts

$$Z_9^2 + (M_9 + V_9) Z_9 \sin(Z_9 T_9) - (U_9 M_9 + Y_9) + (M_9(U_9 + V_9) - X_9)$$

$$\mathrm{X}\cos(Z_9 T_9) - 2M_9 V_9 \cos^2(Z_9 T_9) + M_9 V_9 = 0$$

$$(M_9 + U_9) Z_9 - (M_9 + V_9) Z_9 \cos(Z_9 T_9) + (M_9(U_9 + V_9) - X)$$

$$\mathrm{X}\sin((Z_9 T_9) - M_9 V_9 \sin(Z_9 T_9) = 0$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 = 1 - \cos Z_9 T_9$and $A_9 = U_9 - V_9$we obtain

$$4 M_9^2 V_9^2 W_9^4 + 2 M_9 V_9 W_9^3[(M_9 - V_9)^2 + A_9(M_9 + V_9) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2[M_9 - V_9)^2(2 A_9 M_9 + Y_9 - X_9) + (2A_9 M_9 - X_9)^2 - A_9(M_9 - V_9)(2A_9 M_9 - X_9)$$

$$+ M_9 V_9 (A_9^2 + 4(X_9 + Y_9))]$$

$$+ W_9[A_9^2 + 2X_9)(2 A_9 M_9 - X_9) + 2A_9 Y_9(M_9 + V_9) - M_9 A_9^3] + A_9^2(X_9 M_9 - Y_9) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9(2, U_9 V_9) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting $\cos(Z_9 T_9) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MDDULE NUMBERED SEVEN**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability

criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

## Morphological And Epistemological And Syllogistic Models For Which Nunny's Model Could Be Applied And Has Been Applied:

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production**

**Charge Asymmetry, $\Delta Y \equiv Yt - Yt^{-}$, Reflects The Asymmetry In $Tt^{-}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles

ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic

quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation

between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to

nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of

Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first

impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of

another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease. Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation

of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces

always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the

thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of

thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity

and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI

Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs” Outcomes—

And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities

Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR}$$

$> 0.$ (2)

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

## DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t-T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\big(R(t)\big)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and haimovici model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

**TIME DELAY**

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Limitations of Predator prey analysis in immense. When we realize that world itself is Debit-Credit world and sometime credit-Credit world (See an earlier paper entitled the same nomeneclature.

Opposite limit of absolute stability. The result is a specific The Doctrine of Vibration awareness of an object of the form: 'here this is such.' While individual appearances do not lose their separate identity {svarupabheda) when they rest on a common basis, even so the particular object which appears according to its own characteristics (svalak$ana) is an individual reality in its own right. It is a different kind of appearance characterised by its association with the appearance of the specific location and time in which it is made manifest. The form of our experience is thus 'I now see this here'. But when we perceive each particular constituent appearance separately, each assumes a separate fixed function. Abhinava cites the following colourful example to illustrate how the various combinations of appearances account for the variety of experience: Thus even though the appearance of the beloved may manifest externally, it is as if far away in the absence of another appearance, namely, that of 'embracing'. So when the [appearing of the beloved] is associated with another appearance [namely that of 'far away'] the power (arthakriya) it formerly had of giving pleasure appears as its contrary. The form our experience assumes depends, not only on the nature of the object perceived, but also on personal factors entirely peculiar to ourselves. This theory explains this in two ways. In one sense, the object remains the same, but one or other of its constituent appearances

comes to the fore according to the inclinations of the perceiver. From another point of view, we can say that the perceived object is different for each perceiver according to the difference in the prominent appearance manifest to him. Abhinava, citing as an example a golden jar, illustrates how the same object appears differently to different perceivers according to the use they wish to make of it and to their state of mind: When a person who is depressed and feels that there is nothing [of value for him in the world] sees the jar, he merely perceives the appearance 'exists' [in the form of the awareness that] 'it is\ He is not conscious of any other [of its constituent appearances] at all. An individual who desires to fetch water [perceives] the appearance 'jar'. The man who simply wants something that can be taken somewhere and then brought back [perceives] the appearance 'thing'. The man who desires money [perceives] the appearance 'gold'. The man who desires a pleasing object [perceives] the appearance 'brightness' while he who wants something solid sees the appearance 'hardness'. These 'atomic events' or appearances emerge from the pure subject's consciousness and combine together to form a total event at each moment. Integral Monism of Kashmiri Saivism Daily life (vyavahara) goes on by virtue of this ever renewed flux of appearances. They are connected together and work towards a single unified experience because they appear within the field of consciousness of the universal subject. The aggregate of appearances arises in the [supreme] subject as do [sprouts in] a rice field. Even though each sprout germinates from its own seed, they are perceived as a collective whole. Appearances rest in this way within the universal subject. 'External- ity' is itself another appearance; it arises from a distinction between appearances and the individual subject. So, although all manifestation always occurs within the subject, it appears to be external due to the power of Maya which separates the individual subject from his object. This split must occur for daily life to be possible. Only externally manifest appearances can perform their functions; when they are merged within the subject and at one with him, they cannot do so. Daily life proceeds on the basis of the operation and withdrawal of the conditions necessary for fruitful action to be possible. Appearance in this sense represents the actualisation of a potential hidden in consciousness made possible by virtue of its dynamic, Spanda

nature which is both the flow from inner to outer and back as well as the power that impels it. The emergence from, and submergence into, pure consciousness of each individual appearance is a particular pulsation (visesaspanda) of differentiated awareness. Together these individual pulsations constitute the universal pulse (samanyaspandd) of cosmic creation and destruction. Thus, every single thing in this way forms a part of the radiant vibration {sphuratta, sphurana) of the light of absolute consciousness. Light and Awareness: The Two Aspects of Consciousness Absolute consciousness understood as the unchanging ontological ground of all appearing is termed 'Prakasa'. As the creative awareness of its own Being, the absolute is called ' Vimarsa '. Prakasa and Vimarsa — the Divine Light of consciousness and the reflective awareness this Light has of its own nature — together constitute the all-embracing fullness (purnata) of consciousness. The Recognition (pratyabhijfia) school of Kashmiri Saivism develops this concept of the absolute which finds its fullest expression in Utpaladeva's Stanzas on the Recognition of God. Even though neither of these two key terms appear in the Stanzas on Vibration or the Aphorisms of Siva, they recur frequently in their commentaries. Thus, although the original formulation of the Doctrine of Vibration differs from the theology of Recognition in this respect, it was extended in the course of its development to accommodate this concept of the absolute as well. This was possible, and quite justified, insofar as the absolute understood in Pratyabhijfia terms does not, as we shall see, differ essentially from that of the Spanda school. We can, as Kashmiri Saivites themselves have done, explain one in terms of the other. The Doctrine of Vibration Prakasa: The Light of Consciousness Prakasa is the pure luminosity' (bhdna) or 'self-showing' that constitutes the essence and ultimate identity (atman) of phenomena. That things appear at all is due to the light of consciousness, and their appearing (avabhasana) is itself this Light which bestows on all things their evident, manifest nature. Established in the light of consciousness everything appears there according to its own specific nature (svabhava). Anything that supposedly does not rest in this Light is as unreal as a sky-flower. 3 Thus, according to Rajanaka Rama, unlike the light of the sun, or any other light, this Light not only makes all things apparent, it is also their ultimate source. 4 Full of its

divine vibration the Light makes all things manifest and withdraws them into itself. This supra- temporal activity characterises it most specifically; devoid of it, it would be no better than an inert physical phenomenon. At the same time, this light is the conjunction (slesa) or oneness (aikdtmya) of its countless manifest forms, 6 and the collective whole (Sampina'ana) of all the categories of existence. The universe is nothing but the shining of the Light within itself. It is the radiant vibration (sphuratta) of this Light, the state (avast hand) in which consciousness becomes manifest. Although the Light shines as all things at all times and hence also makes their diversity manifest, 9 penetrating each object individually as well as collectively, it is not totally 'merged' (magna) or identified with the object so as to suffer any division within itself. Our experience of any object is of the form: T see this': it is not itself an object, but the manifest form the object assumes as a luminous principle of experience. The Light is ever revealed and can never be obscured; objectivity can never cast a shadow on the light of consciousness. The Stanzas on Vibration declare: That in which all this creation is established and from whence it arises is nowhere obstructed because it is unconditioned by [its very] nature. This Light is the highest reality (Paramartha). lity with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**DELAY DUE TO CONFEDERATION DEBIT TYPE MATURATION CREDIT TYPE CONSOLIDATION**

**MATURATION**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \quad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9} + \frac{N_9 dG_9}{dR_9} - \frac{dB_9}{dR_9} > \frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9} + \frac{dM_9}{dR_9}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened**

**eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production Charge Asymmetry, $\Delta Y \equiv Yt - Y\bar{t}$, Reflects The Asymmetry In $T\bar{t}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton

particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results

in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely

dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact**

**most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the**

**thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is

nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the

world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease.

Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition

and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the

second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular

chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon

be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining

Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been

Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing

Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

**ABSOLUTE STABILITY**

$$L_9 = \frac{N_9 G_9}{2M_9} \frac{d(9nF_9)}{dR_9} (\geq 0) \quad (9a)$$

$$I_9 = -\frac{N_9 G_9}{2M_9} \frac{d(9nM_9}{dR_9} (\geq 0) \quad (9b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9 > L_9 - I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_9 > L_9 - I_9,$ when $P_9^2/4 > I$ (8a)

$P_9 > L_9 - I_9 + \frac{(I_9 - P_9^2}{(L_9 - I)}$ when $P_9^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to

reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

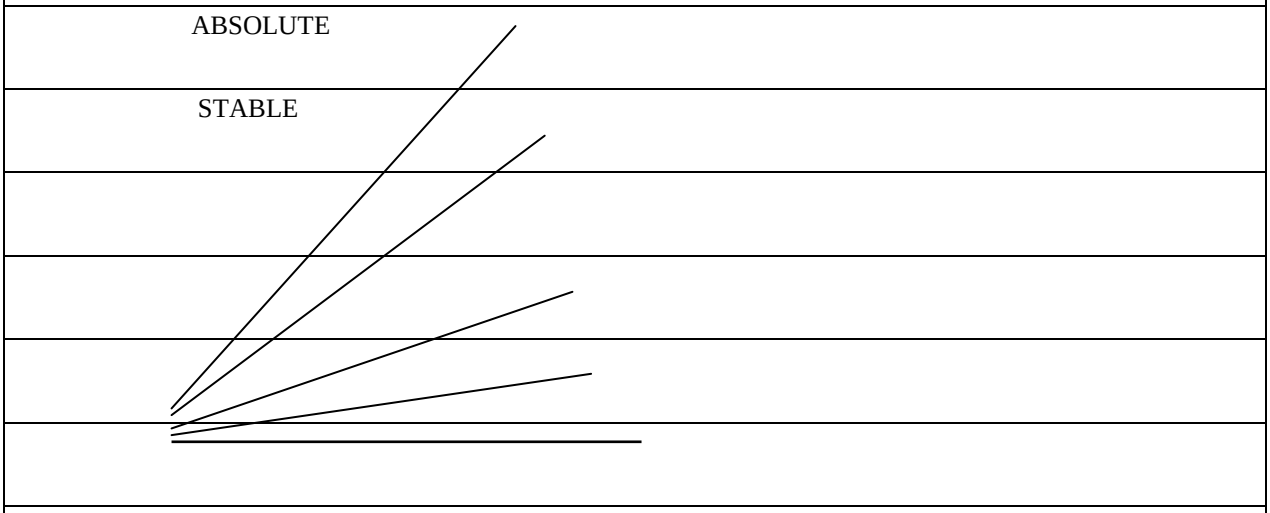



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori

reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

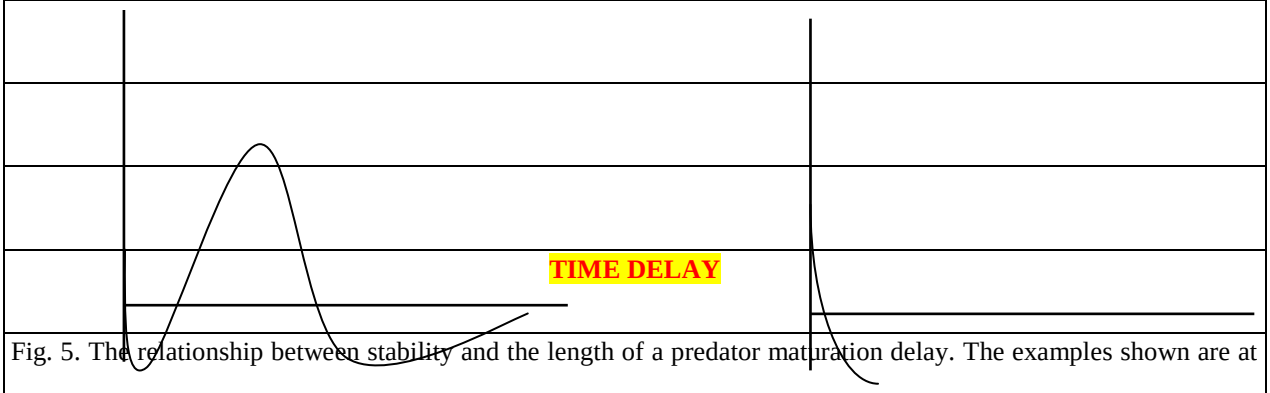



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

## DELAY IN BOTH PREDATOR AND PREY VIS A VIS DEBIT TYPES AND CREDIT TYPES

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t - T_9)F_9\big(R_9(t - T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t - T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) and the effective prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting

$dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

Systems to which model has been applied also include: **Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI:

http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors.

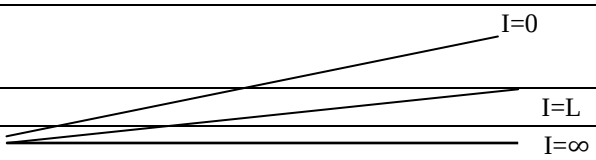
 

Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, ,**

**Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)**

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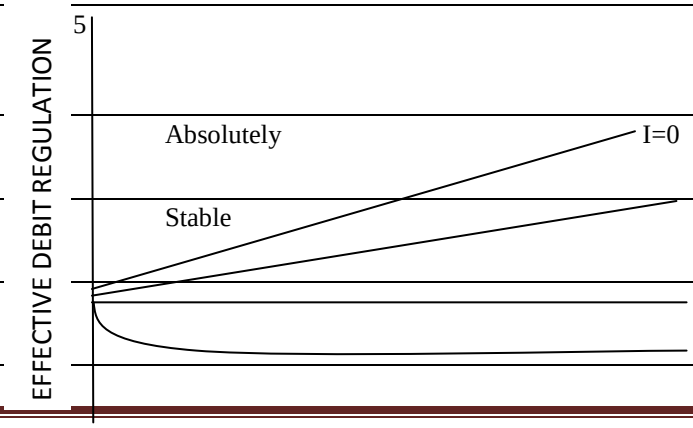



DELAYED DEBIT EFFECT

FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.



-50 10

**DELAYED CREDIT EFFECT (LACK OF BALANCE) CREDIT NON CUMSUMMATION:"**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

The following treatise bears ample testimony and infallible observatory to the fact that how Nunny's model can be applied to various cosmological, general relativity and economic problems.

## ONTOLOGICAL AND EPISTEMOLOGICAL FEATURES OF THE STUDY: OF NUNNY's MODEL

When they are known to be as they are, in the fullest sense of their existence, their phenomenal nature disappears leaving the ground of Being naked and accessible. This approach was validated by a critique of experience. The Vedanta established that space, time and the other primary categories of our daily experience can have no absolute existence. It was therefore necessary to make a distinction between relative truths — that accepted by the precritical common man — and an absolute truth discovered at a higher level of consciousness. The Saiva absolutist 17 rejects any theory that maintains that the universe is less than real. From his point of view a doctrine of two truths, one absolute and the other relative, endangers the very foundation of monism. The Kashmiri Saiva approach is integral: everything is given a place in the economy of the whole. It is equally wrong to say that reality is either one or diverse. Those who do so fail to grasp the true nature of things which is neither as well as both. "We do not" says Abhinavagupta, "base our contention that [reality] is one because of the contradictions inherent in saying that it is dual. It is your approach {paksa) that accepts this [method] While, if [duality and oneness] were in fact [to contradict each other], they would clearly be two [distinct realities]." The Vedantin, who maintains that non-duality is the true nature of the absolute by rejecting duality as only provisionally real, is ultimately landed in a dualism between the real and illusory by the foolishness of his own excessive sophistry (vacafadurvidya). Oneness is better understood as the coextensive unity (ekarasa) of both duality and unity. They are equally expressions of the

absolute. Gopinath Kaviraj says: According to Sarikara, Brahman is truth and Maya is inexplicable (anirvacaniya). Hence the [Advaitin's] endeavour to demonstrate the superiority of Advaita philosophy is turned against his own system. It tarnishes the picture of its philosophical perfection and profundity. He cannot accept Maya to be a reality, therefore his non-dualism is exclusive. The whole system is based on renunciation and elimination and thus is not all-embracing.... By accepting Maya to be Brahman (brahmamayi), eternal (nitya) and real (satyarupa), Brahman and Maya [in the Tantra] become one and coextensive. **The Doctrine of Vibration: An Analysis of the Doctrines and Practices of Kashmir Shaivism Mark S. G. Dyczkowski. Concept of being from Shaivite point of view:** Solutional behaviour and stability analysis of such systems is of paramount importance, notwithstanding the assumptions made which are not farfetched, in that it throws light on the fact whether they could be candidates for yet unsolved problems, providing Authetic ative determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication for the wide amplitudinal ramificatory usage of the thesis propounded. The world is now the other person's world, a foreign world that no longer comes from the self, but from the other. The other person is a "threat to the order and arrangement of your whole world...Your world is suddenly haunted by the Other's values, over which you have no control"(Satre: being and Nothingness) Sartre states that many relationships are created by people's attraction not to another person, but rather how that person makes them feel about themselves by how they look at them. This is a state of emotional alienation whereby a person avoids experiencing **their subjectivity by identifying themselves with "the look" of the other**. The consequence is **conflict**. In order to maintain the person's own being, the person must control the other, but must also control the **freedom of the other "as freedom**". These relationships are a profound manifestation of "bad faith" as for-itself is replaced with the other's freedom. The purpose of either participant is not to exist, but to maintain the other participant's looking at them. This system is often mistakenly called "love", but it is, in fact, nothing more than emotional alienation and denial of freedom. In

essence we are always simulating and selling ourselves at the cost of our freedom. Both as a script writer and participant consciousness, human beings experienceâ complete subjective feelings and sometimes absolute subjective feelings, which if remains unchecked might lead to pathological ramifications. (Italics mine) Perspicuous forbearance, Sophisticated seasoning, Partcipational observation, Background actuality, Existential worldliness, Existential strife, Predicational anteriorities, Character consonance, Ontological consonance, Primordial exactitude, Phenomenological correlates, Accolytish representation, Atrophied asseveration. Anamensial alienisms, Anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orinetatilities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state form the bastion, pillar, post, stylobate and sentinel of the classification scheme and doxa. appropriate and diagnostic, differentiating-discriminating, distinctive-distinguishing, emblematic- especial, essentially exclusive, fixed, idiosyncratic-inborn, inbred, indicative and individual, individualistic-inherent, individualizing and ingrained, inherent, innate, local, marked, native, normal, original, particularly- peculiar, personal, private, proper, regular, representative-singular, special, specific, symbolic, symptomatic- unique these form the classification and stratification measure that is suitable notwithstanding the constancy in many a case, as in gravity, which depends on the mass and distance for two given objects**.** Like Kant and Bergson, Deleuze considers traditional notions of space and time as unifying forms imposed by the subject. He therefore concludes that pure difference is non-spatio-temporal; it is an idea, what Deleuze calls "the virtual". (The coinage refers to Proust's definition of what is constant in both the past and the present: "real without being actual, ideal without being abstract.") Proust, Le Temps Retrouvé, ch. III: see the fourth line from the bottom of this page, or, in English translation, the thirteenth paragraph here: "I began to discover the cause by comparing those varying happy impressions which had the common quality of being felt simultaneously at the actual moment and at a distance in time, because of which common quality the noise of the spoon upon the plate, the unevenness of the paving-stones, the taste of the madeleine, imposed the past

upon the present and made me hesitate as to which time I was existing in. Of a truth, the being within me which sensed this impression, sensed what it had in common in former days and now, sensed its extra-temporal character, a being which only appeared when through the medium of the identity of present and past, it found itself in the only setting in which it could exist and enjoy the essence of things, that is, outside Time. [...] **Nothing but a moment of the past Much more perhaps; something which being common to the past and the present, is more essential than both. [...] a marvellous expedient of nature had caused a sensation to flash to me—sound of a spoon and of a hammer, uneven paving-stones**—simultaneously in the past which permitted my imagination to grasp it and in the present in which the shock to my senses caused by the noise had effected a contact between the dreams of the imagination and that of which they are habitually deprived, namely, the idea of existence—and thanks to that stratagem had permitted that being within me to secure, to isolate and to render static for the duration of a lightning flash that which it can never wholly grasp, a fraction of Time in its pure essence. When, with such a shudder of happiness, I heard the sound common, at once, to the spoon touching the plate, to the hammer striking the wheel, to the unevenness of the paving-stones in the courtyard of the Guermantes' mansion and the Baptistry of St. Mark's, it was because that being within me can only be nourished on the essence of things and finds in them alone its subsistence and its delight. It languishes in the observation by the senses of the present sterilised by the intelligence awaiting a future constructed by the will out of fragments of the past and the present from which it removes still more reality, keeping that only which serves the narrow human aim of utilitarian purposes. But let a sound, a scent already heard and breathed in the past be heard and breathed anew, simultaneously in the present and in the past, real without being actual, ideal without being abstract, then instantly the permanent and characteristic essence hidden in things is freed and our true being which has for long seemed dead but was not so in other ways awakes and revives, thanks to this celestial nourishment "While Deleuze's virtual ideas superficially resemble Plato's forms and Kant's ideas of pure reason, they are not originals or models, nor do they transcend possible experience; instead

they are the conditions of actual experience, the internal difference in itself. "The concept they [the conditions] form is identical to its object."[Desert Islands, p. 36.] A Deleuzean idea or concept of difference is therefore not a wraith-like abstraction of an experienced thing, it is a real system of differential relations that creates actual spaces, times, and sensations. [See "The Method of Dramatization" in Desert Islands, and "Actual and Virtual" in Dialogues. In fact it not as if the style and substance of science and philosophers have been entirely different. look at the conformality and consonance in the Gaussian Theory and Deleuze's statements on identity and the loss of it in the following passage…….That is, identity is a continuous process that finds its limit in exhaustion. As Deleuze will write, "…the logical relation of causality is inseparable from a physical process of signalling, without which it would not be translated into action. By 'signal' we mean a system with orders of disparate size, endowed with elements of dissymmetry; by 'sign' we mean what happens within such a system, what flashes across the intervals when a communication takes place between disparates. The sign is indeed an effect, but an effect with two aspects: in one of these it expresses, qua sign, the productive dissymmetry; in the other it tends to cancel it" (dr, 20). An actualized entity just is such a signal-sign system, remaining in perpetual communication with the virtual multiplicity it actualizes, constituting its own elements and identity. Such is Deleuze's account of what Badiou refers to as the operation of the "count-for-one"… an operation that is immanent to being itself and energetic in character. yet how are we to discern the local nature of ontological situations, the manner in which there is no global, overarching global situation in which all local situations are embedded as if parts of a whole? in part the above passage already responds to this question. if series must be brought into relation to communicate in order for causality to occur, and if the resonance of series is always a matter of chance, of a throw of the dice, like Lacan's objet a where no series enjoys primacy of model over copy, then all we have are local situations without a total situation. Unlike Heidegger's analysis of Dasein where Dasein is always being-in-the-world, Deleuze's simulacra or actualizations are always divergent, enjoying only local relations. Heidegger argues that in order to

understand the being of Dasein we must understand the manner in which Dasein is a being-in-the-world. By contrast, Deleuze's concept of multiplicity provides us with the principle of a new structuralism, of a new local ontology, that allows us to understand the immanent organization of a multiplicity without referring it to an embedding global space. as DeLanda so beautifully puts it in relation to the mathematician gauss, "…when gauss began to tap into these differential resources, a curved two-dimension surface was studied using the old Cartesian method: the surface was embedded in three-dimensional space complete with its own fixed set of axes; then, using those axes, coordinates would be assigned to every point of the surface; finally, the geometric links between points determining the form of the surface would be expressed as algebraic relations between the numbers. but gauss realized that the calculus, focusing as it does on infinitesimal points on the surface itself (that is, operating entirely with local information), allowed the study of the surface without any reference to a global embedding space. basically, gauss developed a method to implant the coordinate axes on the surface itself (that is, a method to implant the coordinate axes on the surface itself (that is, a method of 'coordinatizing' the surface) and, once points had been so translated into numbers, to use differential (not algebraic) equations to characterize their relations. as the mathematician and historian Morris Kline observes, by getting rid of the global embedding space and dealing with the surface through its own local properties, 'gauss advanced the totally new concept that a surface is a space in itself'" (intensive science and virtual philosophy, 11-12).according to the old Cartesian method, we can only outline the properties of a space by relativizing it to a global space in terms of which it is then mapped. by contrast, gauss is able to explore a space in terms of its intrinsic metric and organization as a local space, without referencing it to a whole of which it is conceptualized as a part. **The Image of Thought** "Image of thought "permeates both popular and philosophical discourse. According to this image, thinking naturally gravitates towards truth. Thought is divided easily into categories of **truth and error**. The model for thought comes from the educational institution, in which a master sets a problem and the pupil produces a solution which is either true or false. This image of the subject

supposes that there are different faculties, each of which ideally grasps the particular domain of reality to which it is most suited. In philosophy, this conception results in discourses predicated on the argument that "**Everybody knows..."** the truth of some basic idea. Descartes, for example, appeals to the idea that everyone can at least think and therefore exists. Deleuze points out that philosophy of this type attempts to eliminate all objective presuppositions while maintaining subjective ones. Deleuze maintains, with Artaud, that real thinking is one of the most difficult challenges there is. Thinking requires a confrontation with stupidity, the state of being formlessly human without engaging any real problems. One discovers that the real path to truth is through the production of sense: the creation of a texture for thought that relates it to its object. Sense is the membrane that relates thought to its other. Accordingly, learning is not the memorization of facts but the coordination of thought with a reality. "**As a result, 'learning' always takes place in and through the unconscious, thereby establishing the bond of a profound complicity between nature and mind".** Deleuze's alternate image of thought is based on difference, which creates a dynamism that traverses individual faculties and conceptions. This thought is fundamentally energetic and signifying: if it produces propositions, these are wholly secondary to its development. Philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; even tum tan tum;it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the vorticity of being merges with the positive use of the disjunctive synthesis, and

this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoin der for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for every hum of voices or a single silence for all the silences; a single vacuum for all the vacuousness consciousness should not be said without conjuring if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. The attributable to all states of bodies and states of affairs and the expressible of every proposition So univocity of consciousness means the identity of thé noematic attribute and that which is expressed linguistically and sense-fully. Univovity.We mean that it does not allow consciousness to be subsisting in aquasi state and but expresses in all pervading reality. (Reference: LOS: Deleuze) Deleuze sometimes spoke of getting up behind an author and "creating a monster", as if his reading methodology somehow consisted of a distortion of the philosopher's or author's thought. Certainly there is a case to be made for this with regard to the process of deterritorialization, where something is wrested from a territory and deterritorialized upon a new territory, like the animal paw that is deterritorialized from the earth and reterritorialize on the branch. however, it seems to me that the more interesting aspect of Deleuze's approach to other philosophers and art is not so much his "monstrous creations" (i seldom find them particularly monstrous), but rather their Gaussian or Riemannian style, where he explores them in terms of their own internal organization and metric, without reducing them to something alien such as history, society, biology (reductivism), or the signifier. What we find in Deleuze's approach to phenomena is a Gaussian technique. for instance, take Deleuze's books on cinema, his book on Francis bacon, or his study of sacher-masoch. In the first instance, Deleuze's carefully separates cinema from narrative and the signifier, studying it in terms of its specific organization pertaining to the production of images. in the case of bacon, Deleuze doesn't look for an underlying narrative or "meaning", but instead studies the manner in which bacon composes and

organizes his images and lines, both in terms of their production and actuality, so as to liberate a logic of sensation. finally, in approaching sacher-masoch, Deleuze allows Sacher-Masoch's novels to speak for themselves in terms of their desire and relation to pain, stalwartly refusing to reduce masochism to the complement of sadism. in each case, we have a local exploration of a "space" of multiplicities that is extremely precise. Sean Bowden the priority of events: **Deleuze's logic of sense sean Bowden, the priority of events: Deleuze's logic of sense, Edinburgh university press, 2011, 296pp., $40.00 (hbk), isbn 9780748643646 Reviewed by David Scott, Coppin state university and Gaussian spaces and multiplicities a note on Gaussian spaces and multiplicities posted by larval subjects under uncategorized.** What you "see" is not always what you "get." Many people mistakenly take their own visions literally without "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in the physical sense but in a psychological sense. The archetypes of the collective unconscious are arrayed behind the scenes of current worldwide conditions, of crisis and confusion. They mirror our own states back at us, whether we perceive them as such or interpret them plausibly or not. The noise of ordinary consciousness and beliefs drowns out the signal. Unconsciousness is the background of our ordinary awareness. Our organism is very much at the center of such effects. The organismic source is our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is (=) a meme-Scape. Pre-conceived concepts vie with (e&eb) structures, concepts with images. Like scientists who ignore (e) assumed truths, we leapfrog over (e&eb) our beliefs and personality deficits, claiming (eb) idiosyncratic imagination is (=) literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves (e) us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our

experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images. It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is (=) both. We inherit (e&eb) the structures that make our experience (=) what it is. But the structure itself is "empty," and each human culture "fills" it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about (e&eb) our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from

criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns,

nomenclature, and internal coherence of meaning and representation. n a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar some statements in Vedanta, it is not superfluous or redundant. It is evolutionary. Authetication determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication of the evolutionary process towards understanding is ever present in all, its thematic and discursive form. it is the relentless holding on to a theory this is despicable albeit the theory is not. Look at this paragraph. It is the 'I am the body/mind' belief that gives rise to the 'I am not (e) the world' belief. These two beliefs are (=) co-created."Consciousness projects (e&eb) the appearance of the mind, body and world by taking the shape of thinking, sensing and perceiving." "Attention is Consciousness with an object. When the object vanishes, attention simply remains what it always is, Consciousness.""There is (=) no purpose to (e) meditation. The purpose is already accomplished." "Everything that is experienced is experienced by, through, in and as Consciousness." "The seen cannot be separated from seeing and seeing cannot be separated from (e) Consciousness.""The Reality of any experience is not hidden in the appearance; it is expressed by the appearance." "Once we see that everything is Consciousness... Maya still dances, but it is a dance of love not seduction." There are some truly excellent sections

such as a long one taking, as a starting point, the following quotation from the painter Cézanne: "Everything vanishes, falls apart, doesn't it? Nature is always the same but nothing in her that appears to us lasts. Our art must render the thrill of her performance, along with her elements, the appearance of all her changes. It must give us a taste of her Eternity." Rupert begins: "That statement must be one of the clearest and most profound expressions of the nature and purpose of art in our era." And he goes on for some 14 pages to elaborate on this claim, examining the nature of the 'elements' Existence, appearance and Consciousness and their relationship

Some of the following concepts are modeled; some are not. For those which are not modelled kindly see one of the papers in the series. This note has the axiomatic predication to give a holistic view of the study. Here we talk of the characterstics of systems which satisfy the condition of cosmological constant or any other for that matter or variable or axiomatic predication. . There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit,

facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified

from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. Towards the end of circumvention of typing hundreds of superscripts and subscripts which would be a sardine squash and the concomitant operational difficulties, model is presented in piece meal of nine modules. Logarithms are to be taken in respect of those which incorporate more than one variable in bra-ket. Values of log (ab) and log (a+b) are readily available. Anti logarithm shall be taken at the end to predict, project, prognosticate the value of the variable. This is true for tensors, vectors and other variables too. Affirmational assertion, and explanation justification, statement of vindication, annotational commentary, and explicational glossary for each and every system changes and physical interpretation of results is one thing that is to be with earnest endeavour and feverish and febrile expediency. Editions never ever mean the same and identicality of parameters and this has been explicitly and unmistakably stated in the model apriori itself. Akin and analogous, cognate and concurrent, correspondent and congruent, comparable and complementary, synonymous and duple, tantamount and agnate, commensurate and correlative representation is only to highlight the importance and subterfuge the replication of work. It is my fervent solicitation to kindly bear with me for any lapses, notwithstanding the orchestrated efforts for a paper without any minor errors. Postulation predication, conclusive presumption, differential presuppositions, underscored decidedly axiomatic statement of the statement, equation form the bye word or the watch word in the aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement-argument, articulation, assertion, asseveration, choice of words, commentary and communication, declaration and definition, delivery-diction, elucidation-emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation tone and tenor of stratification. Now, how do find the reaction of systems to these singularities. You do the same thing a boss does for you. "Problematize" the events and see how you behave. I will resort to "pressure tactics". "intimidation of deriding report", or "cut in the increment" to make you undergo trials, travails and tribulations. I am happy to see if you improve your work; but may or may not be sad if you succumb to it and hang yourself! We do the same thing with systems. systems show conducive response, felicitous reciprocation or behave erratically with inner roil, eponymous radicality without and with blitzy conviction say like a solipsist nature of bellicose and blustering particles, or for that matter coruscation, trepidiational motion in fluid flows, or seemingly perfidious incendiaries in gormandizing fellow elementary particles, abnormal ebullitions, surcharges calumniations and unwarranted(you think so but the system idoes not!) unrighteous fulminations. Amount of emotion or affect that is appropriate or expected in an given form of interaction. Particular individuals and diffuse obligations (see c and d) are associated with affectivity, whereas contacts with many individuals (universalistic) in a bureaucracy may be devoid of emotion and characterized by affective neutrality. Affective neutrality may refer to self discipline and the deferment of gratification (eg Weber's spirit of capitalism). In contrast, affectivity may be associated with expressing emotions. Adams and Sydie also refer to affective neutrality being associated with ego control (p. 15). (Parsons: Wikipedia). So the point that is made here is "like we problematize the "events" to understand the human behaviour we have to "problematize" the events of systems to understand their behaviour. This statement is made in connection to the fact that there shall be creation or destruction of particles or complete obliteration of the system (blackhole evaporation) or obfuscation of results. Some systems are like "inside traders" they will not put signature atoll! How do you find they did it! Anyway, there are possibilities of a CIA finding out as they recently did! So we can do the same thing witn systems to. This is accentuation, corroboration, fortification, .fomentatory note to explain the various coefficients we have used in the model as also the dissipations called door. In the bank example we have clarified that various systems are individually conservative, and their conservativeness extends holisticallytoo.that one law is universal does not mean

there is complete adjudication of nonexistence of totality or global or holistic figure. Total always exists and "individual" systems always exists, if we donot bring Kant in to picture! For the time being let us not! Equations would become more ensorcelled and frenzied..... philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". In consciousness research, two rival sets of theories can be recognized: (A) Scientific material interpretations of consciousness are based on axioms that view consciousness in the context of highly advanced intentional processing of information in which subject-object relations evolve, and (B) humanistic interpretations of consciousness are based on axioms that view consciousness in the context of, say, "centered pulsations" that enable a conscious agent to act from his or her center of awareness. In this paper I will argue for the selection of axioms that favor humanistic interpretations of consciousness. The Pursuit of Autonomy Interdisciplinary Observations to Human Consciousness: Wautischer, Helmut (1993) the Pursuit of Autonomy. Interdisciplinary Observations to Human Consciousness: Social Neuroscience Bulletin, 6 (4). pp. 52-56.Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality; **CONSCIOUSNESS AND ITS UNIQUENESS June 12, 2012 at 8:40pm (See Deleuze Logic of sense, Wikipedia and Stanford encyclopedia for more details)** Atrocious contrivance and device, gratificational primogeniture, calamitous dodge and expedient, depraved, destructive, disastrous, execrable gambit and gimmick, great solace and succor, iniquitous, injurious, loathsome, low, maleficent, malevolent machination and maneuver, promethaleon of candor, frankness, honesty, honor, ingenuousness, innocence, openness, reality, sincerity, truthfulness, spiteful, stinking, ugly, unpleasant, unpropitious play and ploy, progenitor of jurisprudence and circumspection, racket and ruse, savvy, scam, stratagem, subterfuge, tactic and wile, forthrightness, honesty, truthfulness embodied and personified is the structural predisposition and dispensation of the complex. Nobody need have to take the Brahman-Anti Brahman agency to be a recalcitrant, repugnant, and refractory proposition. God gives and also takes away. That is the point made. ‘

And of course lesser gods, and God doth follow them and they follow God. That is why they are lesser gods. And wield the power on lesser mortals like us. When we mean gratification and deprivation we mean celestial Brahman Anti Brahman and Terrestial Brahman and Anti Brahman agencies producing such happiness or sadness by any means such as (blending- amalgamation, adulterant-adulteration, amalgam-amalgamation, blend- combination,

composite- compound, debasement-denaturant, fusion-hybrid, intermixture-reduction, amalgam- mixture, composite- compound, fusion-mishmash, aggregate- alloy, amalgamation-blend, commixture-composite, composition-compost, conglomerate-fusion, goulash-medley, mishmash-stew, blend-coalition, commixture-compound, federation-heating, integration-junction, melting-merger, synthesis-unification, assortment- combination, adulteration-alloy, assimilation-association, batter-blend, brewed combine, composite compound, concoction confection, conglomeration cross of money, orientation(dis), penury creation, murder, mayhem plunder, pillage, apocalypse, Armageddon in various forms and permutation and combination of Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness: with each action accounted and showing balance in the nature's general ledger as also actionlessness:e). All these are recorded in Nature's general ledger. Neuron DNA encapsulates all the actions passions, actions, interactions, transactions or the lack of it at various levels of individual, collective and cosmic levels. In celestial Brahman Anti Brahman ledger there is cessation of dualities. One more point to be made is if 5=3+2, then we can say 5 gormandizes 3 or 5 gobbles up 2. Quintessentially such a statement arises out of the difference of LHS and RHS term wise. It must be remembered despite all theoretical abstractions and generalizations be it proton mass or Higgs Boson, or topology of a space, it is notwithstanding the fundamentality of nature is just a number. Just everything in Physics and mathematics is just a number. Let us assume I buy rice in Bangalore (District) worth 5k.g. I can always divide 5k.g.s of rice in categories of Bangalore

(South), Bangalore (west), Bangalore (East) and Bangalore (South). On similar basis, there are billions of neutrons, protons and concomitant interactions which would satisfy the given numerical value. Example of gravity above makes things clear. Everything atleast is a function of time and someplace or some other coordinates, subordinates or superordinates. Mass of proton is 2.47 does not mean total mass of all protons is 2.47. It is this argument that is used in the classification scheme. By the very nature, self the unmonitoring but the witness consciousness agency identifies itself with something, be it a profession, an actor on the screen, when it does not have a identity. Such examples are legion, a Charlie Chaplin in "The Kid" or Raj kapoor in "Jagte Raho". These are people who just wanted to live and living without being put under surveillance every moment is the greatest part of their life. This is "Absolute Subjectivity" of Shiva. Such a state is always craving for identification for in pure consciousness all identities are lost. It is this "identitylessness" or having very minimal identity is what pervades most in India and this is due to expansion of individual consciousness. When that happens, be it education, or wealth, the "identity' becomes so well entrenched, they become the exact opposite of "pure consciousness"," impure consciousness" (Shakti). This is a state of "Absolute or relativistic objectivity". There was this case of a lady who went to such an extent of planning, she was sure that whatever she does must work out. When it did not, she lapsed in to "identity crisis". Such thing as this might happen in well developed countries where the competition is extremal and to be achieved at any cost. Realisation that one is not the doer but the witness consciousness is just like watching a movie in which one has acted. It is in this sense "advaita", "non duality' seems to be of cardinal and preeminent importance for the analytical treat ment of "identity' and "identity lessness", both happen in space time. "Multiverses" in this sense is also like a different film of an actor released at different theaters and it is in this sense we state that that every object in space time is a multiverse. You must have totally immersed in doing a "string theory" problem or watching a "film", oblivious of surroundings. Essentially creation (space time) also is what is happening in mind and honed in to you billion times. Identification of anything with a designation or signification is one that leads to

association within spacetime, which would only have pathological consequences and detrimental ramification as stated above or positive implications of "happiness" because of the identity itself. Search or 'show" has to go on! It gives a sense of consummation. a watching pot never boils; donot go to check nature; you are disturbing it; it will never show its true colours bhakti and Virakthi, gratification and deprivation, raga and dwesha is essential to live in space. It is the quantum interest thereof that produces energy to live .gratification and deprivation which form the bastion, pillar, post and stylobate and sentinel of all human actions depends on tams, rajas, and sattva of individual or insentient object. Verily tamas, rajas and satva are the characterstics of space itself, nay the behavioural pattern and attitudinal orientation. There are peaceful places like Vienna or Sringeri and there are violent places like others. Space is warped and woofed by the Akshara, the quantum information of primordial syllable Om. cosmic general ledger is the one vibration of which is brahman; quantum information possessions thereof form the sentinel of all transactions (Gargi to Yajnavalkya: my interpretation)adhocism, nonseqientiality of events, anti sequentiality of events, simultaneity of phenomenon, concurrent happenings , improbability of phenomenon happening itself are characterstics of brahman-anti brahman (Terrestial and celestial) are the cause of non predictability of happenings like in black hole, or in the areas of operation of brahman-anti brahman (Terrestial and celestial). brahman AntiBrahman actions, interactions, decisions, circumspection, adhocism, satisfies sensitive to initial conditions; topologically mixing; dense periodic orbits, the axiomatic predications and postulation alcovishness of chaos theory. Like gravity is a characteristic of space, so is tamas, rajas, sattva, with spaces knotted together by quantum information (for models see below) .accidents are preplanned actions of Terrestial and celestial Brahman AntiBrahman. All actions are perpetrated by them, be it train on fire or death. Belief systems play an important role in the determination of tamas, rajas, and satva state of human beings. There is neither cause nor effect. All actions are performed by Terrestial and celestial Brahman anti Brahman and it is off the record. psychic energy is negative equivalent to –mc^2, so that second law of thermodynamics is satisfied on a holistic basis vis-à-vis material energy

ParaBrahaman is infinite individual, collective and cosmic consciousnesses are infinite (concomitant general ledgers) for they are in the evolution mode with respect to each other in that order. Later emanates from the former as transactions emanate from the conservative state of assets and liabilities. Is not zero true for infinite conservative systems inside space and time? zero is the only consummation, consolidation and conservation. if fifty out of seventy members come out of the room at whatever time you come out, talk about the grand limbo of oblivion and hibernation of lifeless world you are going to and if you are not a celebrity, then it is not coincidence, it is coordination. Nature is playing the same trick on you. That a generalised law is there and you can measure does not make things less good albeit recalcitrant, refractory and rogue particles occur time and again, by and large. Entering in to all living beings through his constant part on the new moon night, he is born therefrom next morning. Absolute ParaBrahaman (zero) projects himself again and again on the screen of individual consciousness which appears of depairing despondent world. Some where the balancing act is done on karma. cosmic general ledger reflects upon itself to keep all transactional ties conservative (Brihadarnkayaka Upanishad: 1.5 14) for models see next paper through its divine power ( i interpret man creates god; it is attribution and ascription of various Shaktis would the formation of god takes place like space and time) the self assumes (eb) the form of a deity man can contemplate and venerate, even though siva, the pure subject, can never in fact be an object of meditation (adhyeya). Until we realise our true identity with Sankara, he is worshipped and conceived to be a reality alien to ourselves. While we are in the realm of creation, he too is a creation or mode or appearing of the absolute, manifest to us in meditation, through his freedom as an eternal, omniscient being. There is no gulf between the created and the uncreated creator: nothing in reality, although an object of knowledge, ceases to be Siva: this is the reason why meditation [on this or that aspect] of reality bestows its fruit. The world of the senses and mind appears to the well awakened (suprabuddha) as a theophany, an eternal revealing of god in his creation. The doctrine of vibration declares that "there is no state in word, meaning or thought, either at the beginning, middle or end, which is not Siva." 59 to utter any word is, in

reality, to intone a sacred formula.every act is a part of Siva's eternal cosmic liturgy, every movement of the body a ritual gesture (mudras), and every thought, god's thought. By what path are you not attainable? What words do not speak of you? in which meditation are you not an object of contemplation? What indeed are you not, o lord? 61 spandashakti, which accounts for the appearing of all things, is also the means by which deity in its many varied forms appears to man. Ksemaraja concludes: the ultimate object of worship of any theistic school differs not from the Spanda principle. The diversity of meditation is due solely to the absolute freedom of Spanda. 62 Sankara is not only the supreme object of devotion; as the static polarity of the absolute, he is the inner reality which holds together its Siva and Sakti outer manifestations. Phenomena are patterns of cognitions projected onto the surface of self-luminous siva-consciousness. There they become apparent, directly revealed to consciousness according to their manifest form. Siva is accordingly symbolized as the ground or surface of awareness, smooth and even like a screen (samabhittitalopama). 64 inscribed on this screen (kutfya) are the countless manifest forms which appear within it rendering it as diverse and beautiful as a fossil ammonite (falagrama). 65 Siva is the sacred ground upon which the cosmic mandalas are drawn, the absolute surface of inscription which bears the mark (cihna) of the universe. Abhinava writes: the variety of this world can only be manifest if the highest lord, who is essentially the pure light of consciousness, exists; just as a surface is necessary for a picture. if external objects were perceived in isolation then, because 'blue' and 'yellow', etc., are self-confined and the perceptions [we have of them] refer to their objects alone and so are insentient, mute and dumb in relation to one another . . . how would it be possible to be aware that an object is variegated? but just as depths and elevations can be represented by lines on a smooth wall, and we perceive [a female figure and think], similarly it is possible to be aware of differences in the variegated (contents of experience) only if all the diverse perceptions are connected together on the one wall of the universal light of consciousness. 66 Siva is the perfect artist who, without need of canvas or brush, paints the world pictures. The instant he imagines it, it appears spontaneously, perfect in every respect. The colours he uses are the varying shades

and gradations of his own Spanda energy and the medium his own consciousness. Intractable disaffect, Mutinous obdurate, Rebellious recalcitrant, Seditious refractory, Insubordinate insurgent, antagonistically antipodean, discrepant incompatible, Contumacious coadjutant, Conflicting converse, Unyielding obdurate, the antagonistic and anachronistic dispositional ties of Anti Brahman are in fact the sine qua non of Brahman only (See Vishwaroopa of Mahavishnu, which contains the polarised positives on RHS and negatives on LHS). The universe is coloured with the dye of its own nature (svabhava) by the power of Siva's consciousness (a//). Rajanaka rama says: homage to him who paints the picture of the three worlds, thereby displaying in full evidence his amazing genius {pratibha)\ to Sambhu who is beautiful with the hundreds of appearances laid out by the brush of his own unique, subtle and pure energy. 68 analogously, at the microcosmic level, all the cognitions and emotions, etc., which make up the individual personality form the outward flow of essentially introverted consciousness. They are specific pulsations (vise$aspanda) or aspects of the universal pulsation (samanyaspanda) of pure t consciousness. at the lower level, within the domain of maya, they represent the play in the fettered soul of the three primary qualities (gunas) or feeling-tones' which permeate to varying degrees his daily experience. These are: 1) sattva — the quality of goodness and luminosity which accompanies blissful experience both aesthetic and spiritual. 2) Rajas — the passion or agitation which oscillates between the extremes of 'light' and 'darkness' and characterises inherently painful experiences. 3) Tamas — the torpor and delusion which accompany states of inertia and ignorance. 69 the liberated soul recognises that these three are the natural and uncreated powers of pure consciousness. for him they are manifest respectively as: 1) Sankara 's power of knowledge (jnana) — the light of consciousness (prakasa); 2) the power of action (kriya) — the reflective awareness of consciousness (vimarsa)\ 3) the power of maya — which does not mean here the world of diversity, but the initial subtle distinction which appears between subject and object in pure consciousness. 70(doctrine of vibration: mark S.G. Dyczkowski) .to "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new

concepts and meanings that could not be included in the old system. at the heart of western metaphysics, for example, derrida finds the opposition between "speech" and "writing." this binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior in derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. The postmodern turn in philosophy: theoretical provocations and normative deficits by Steven best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html. You cannot see beyond consciousness for that is the limit of your storage of information or the basis for understanding further information. Turiya is empty bucket; it is from which prajna the measure of  arises; turiya measures the number of quantum information stores in other three states of waking, dreams, and dreamless deep sleep (Manduka Upanishad; my interpretation for models next paper).let there be path where angels tread: it is a force never born never dies the brahman an anti brahman verily there are parallel universes in (eb) in the ambit of space and time; the society of sinners and the sinned with the grammar of evil; parallel universes contain(e) multiverses in them; all physics, philosophy,  break down(e) in these universes; astronomical gratification ,deprivation, glorification, mortification, projections, axes, dimensions, rotations, foldings, concurrences, motions, shocks, are the axiomatic predications of multiverses and parallel universes; death is when self (the witness consciousness; register of individual general ledger) leaves the body; ostensibly parallel seem to have (e) an objective; parallel universes donot (e) have objectives; their paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common

attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, affection- appreciation, awareness- discernment, emotion-feeling, gut reaction, heart insight, intuition, judgment, keenness, perceptiveness, rationale, sensation, sense, sensitiveness, sensitivity, sentiment, susceptibility, taste, vibes, admiration, aesthetic sense, affection, appraisal, assessment, attraction, awareness, cognizance, commendation, comprehension, enjoyment, esteem, estimation, grasp, high regard, knowledge, liking, love, perceptional realization, recognitional regard, relish, respect, responsiveness, sensibility, sensitiveness, sensitivity, sympathy, understanding, valuation, acknowledgment, airing, baring, , confessional defenselessness, denudation- denunciation, disclosure, display, divulgence-exhibition, hazard, introduction, jeopardy, laying open, liability, manifestation, nakedness, openness, peril, presentation, publicity, revelation, risk, showing, susceptibility , susceptiveness,susceptivity, unfolding, unmasking, unveiling, vulnerability, vulnerableness, action, affection, appreciation and ardor, behavior, capacity, compassion, concern, cultivation and culture, delicacy, discernment and discrimination, emotion, empathy, faculty, fervor, fondness, heat, imagination, impression and intelligence, intensity and intuition, judgment, keenness- palpability, passion, pathos, pity, reaction, refinement, sensibility, sensitivity, sentiment, sentimentality , sharpness, spirit, sympathy, tangibility, taste, tenderness, understanding, warmth, fond memories, hearts and flowers, homesickness, longing, pining, reminiscence, ,schmaltz, sentimentality wistfulness, yearning satiation is only ostensible; their job is to get the job done; their business is to mind other's business ;all laws of physics and finance breakdown; it is antilaw,piscatorial Piratish and bubonic bucaneerishness; there is free transactions between in multiverses; when reflected on the screen of consciousness they appear as extant existential universe. Transactions are astronomical and gratification

deprivation, glorification and mortification incalculable; they have no barriers of region, religion, language which they create for others. theirs objectives are: corruption criminality, debasement debauchery, degradation evil immoral lewdness licentiousness, perversion sinfulness, vice viciousness, vileness wickedness, improbity corruption, atrocity decadence, degeneration degradation, depravity evil, immorality impurity, infamy iniquity, looseness lubricity, perversion profligacy, sinfulness turpitude,vice viciousness, vulgarity wickedness,disgrace bad reputation, abasement abuse, baseness black eye, blemish blur, brand comedown,contempt contumely, corruption culpability, debasement debasing, defamation degradation, derision disbarment, discredit disesteem, disfavor dishonor, disrepute disrespect, humbling humiliationignominy ill repute, infamy ingloriousness, meanness obloquy, odium opprobrium, pollution ,put-down reproach, scandal scorn, slander slight, slur spot, stain stigma, taint tarnish, turpitude venality, enormity Horribleness , abomination atrociousness, atrocity crime and the positive aspectionalities thereof .let there not be any incursions where evil mercenaries in the garb of mendicants thrive with adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatious gainsaying or conformingly pugnacious Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, execrable and goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness . statement for

Mephistophelean mercenaries philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality. Negative energy stabilising wormholes are considered candidates for (er) darkmatter. An analysis of negative energy might be a useful ingredient of the thoughts of travel at speeds more than the velocity of light. Peirce explains that an "imperfect continuum" possesses "topical singularities, or places of lower dimensionality" which break continuity, while "the parts of a perfect continuum have the same dimensionality

as the whole." (CP 4.642) Circular continua without singularities are time and space. (CP 6.210-212) By using math as his lens, Peirce envisioned singularities or lower dimensional holes in the imperfect cosmologic continuum, while maintaining that a probabilistic law of chance exists that entropy decreases **(Wikipedia).** On the side of philosophy**,** quite explicitly, Deleuze provides a critique of Husserl's "sleight-of-hand" change from a static categorical phenomenology to one supposedly more dynamically constitutive of reality. This a clear cut **chagrined circumvention**, **contraventional curbing**, **defeated disgruntlement**, **dissatisfactional downer**, **dragged failure**, **fizzled foil**, **grievance hindrance**, **impedimental irritation**, letdown, nonfulfillment, nonsuccess, obstruction, old one-two, **resent mental setback**, unfulfillment of the earlier theories makes Sartre contend that human existence is a conundrum whereby each of us exists, for as long as we live, within an overall condition of nothingness (nothing-ness)—that ultimately allows for free consciousness. But simultaneously, within our being (in the physical world), we are constrained to make continuous, conscious choices. It is this dichotomy that causes anguish, because choice (subjectivity) represents a limit on freedom within an otherwise unbridled range of thoughts. Subsequently, humans seek to flee our anguish through action-oriented constructs such as escapes, visualizations, or visions (such as dreams) designed to lead us toward some meaningful end, such as necessity, destiny, determinism (God), etc. Thus, in living our lives, we often become unconscious actors—Bourgeois, Feminist, Worker, Party Member, Frenchman, Canadian or American—each doing as we must to fulfill our chosen characters' destinies. However, Sartre contends our conscious choices (leading to often unconscious actions) run counter to our intellectual freedom. Yet we are bound to the conditioned and physical world—in which some form of action is always required. This leads to failed dreams of completion, as Sartre described them, because inevitably we are unable to bridge the void between the purity and spontaneity of thought and all-too constraining action; between the being and the nothingness that inherently coincide in our self. Deleuze insists on the absolute non-resemblance between what conditions (metastable, pre-individual plane of singularities) and what is conditioned

(actualized or the individuated 'thing'). Only then do "the conditions of the true genesis become apparent," Deleuze writes (Logic of Sense 105). Without addressing this issue, Deleuze's borrowings from Simondon, Lautman and structuralism remain untethered. Because of Bowden's decision to completely sidestep Deleuze's overt negotiations with phenomenology, his otherwise sophisticated commentary suggests a pedagogical narrowness. If, the omission of phenomenology strengthens Bowden's focus on the ontological priority of the event over substantial ontologies via the genealogy of figures and movements Deleuze outlines, it also incompletely articulates the problem of the sense-event as Deleuze formulates it. We have lost any idea of whom Deleuze takes to be his primary adversary, against whom he stakes out his position. This worry is most clearly manifested in Bowden's description of Deleuze's ontological prioritizing of events as offering a "transcendental ontology." For Bowden Deleuze's transcendental ontology defines events that "are ontologically prior to substances, 'all the way down'" (82). But Bowden provides a second definition: it is that the priority of events over worldly individuals means that the world of events "is not something external to the conditions of knowledge" (69). **SEAN BOWDEN the Priority of Events: Deleuze's Logic of Sense Sean Bowden, the Priority of Events: Deleuze's Logic of Sense, Edinburgh University Press, 2011, 296pp., $40.00 (hbk), ISBN 9780748643646. Reviewed by David Scott, Coppin State University Weltschmerz, agony, apprehension, blues, depression, dread, mid-life crisis, misgiving, nervousness, uneasiness, angst were expressed with such effulgent words…… aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement, argument, articulational assertion, asseverational choice of words, commentary and communication, declaration and definition, delivery, diction, elucidation, emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation about the human predicament is brought out in Being and Nothingness. Being and Nothingness offers a critique of Sigmund Freud's theory of the unconscious, based on the claim that consciousness is essentially self-conscious. Sartre also argues that**

**Freud's theory of repression is internally flawed. According to Sartre, in his clinical work, Freud encountered patients who seemed to embody a particular kind of paradox—they appeared to both know and not know the same thing. In response, Freud postulated the existence of the unconscious, which contains the "truth" of the traumas underlying the patients' behavior. This "truth" is actively repressed, which is made evident by the patients' resistance to its revelation during analysis. Yet what does the resisting if the patients are unaware of what they are repressing? Sartre finds the answer in what Freud calls the "censor". "The only level on which we can locate the refusal of the subject," Sartre writes, "is that of the censor.** The Holographic Principle, as conceived in current physics, applies to fields, and perhaps, even to more elementary entities, called strings and branes. These more elementary entities remain hypothetical at the time of this writing. There are many layers between the level of fields and that of neurons. We know that neurons are surrounded by a surface phospholipid membrane, which supports an electrochemical process that is fundamentally necessary to human experience. As the information supported by the membrane surface according to the Holographic Principle is proportional to its area, we should expect to find that the neural correlates of consciousness most clearly associated with those parts of the neuron that have the have a high proportion of neuronal surface area, and, in particular, receptive surface area. Dendrites account for an average of about 90% of the neuron is receptive surface area (Wong, 2002), and so would be expected to be the most important neural structures in information processes. We find this to be true **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science (Wikipedia).** Discordant-discrepant, dissonant-hard-line, implacable, incompatible, incongruous- inconsistent, inexorable-inflexible, inharmonious-intransigent, opposed, reluctant, unappeasable, uncompromising, unfriendly some of these expatiations and enumerations might sound, a more generalised interpretations is given by a Vedantin. Look at this arresting passage: The teacher continued to be silent. When addressed a second and third time he said: "I am teaching, but you do not follow. The Self is silence." The undetermined and unthinkable character of the Brahman is

a consequence of the absolute's eternal and immutable nature. To concede the existence of a real universe is, from the Vedantin's point of view, to posit the existence of a reality apart from the Brahman. Nor can we simply identify a real universe with the absolute unless we are prepared to compromise its unchanging, absolute status. The criterion of authenticity is immutability. Reality never changes; only that which is less than real can appear to do so. Reality is constant in the midst of change. What this means essentially is that there is change although nothing changes. This impossible situation is reflected in the ultimate impossibility of change itself. That which does not exist prior to its changing and at the end, after it has changed, must be equally non-existent between these two moments. Although the world of change appears to be real, it cannot be so. Change, according to the Vedantin, presupposes a loss of identity. Reality cannot suffer transformation; if it were to do so, it would become something else and the real would be deprived of its reality. The immortal can never become mortal, nor can the mortal become immortal. The ultimate nature of anything cannot change. Change of any sort is merely apparent (vivarta); the world of change and becoming is a false super- imposition (adhyaropa, adhyasa) on the absolute. In cosmic terms, the mistake (bhranti) consists of the supposition that the real Brahman is the unreal universe and the unreal universe is the real Brahman. In microcosmic terms, it is the mistake of falsely conceiving the body, mind or even one's personality to be the Self. In the same way as the image of a snake is falsely superimposed on a rope, similarly the universe is falsely projected onto the real substratum, the Brahman. Ignorance is not merely a personal lack of knowledge, but a cosmic principle. As such it is called "Maya," the indefinable factor (anirvacaniya) that brings this mistake in identity about. The reality status of this cosmic illusion is also indefinable: on the one hand it is not Brahman, the sole reality; on the other hand it is not absolutely non-existent like a hare's horn or the son of a barren woman. Brahman is the source of world appearances only in the sense of being their unconditioned ground or essential nature. The universe is false not because it has no nature of its own but because it does have one. Just as the illusion of a snake disappears when one sees that it is nothing but a rope, similarly cancellation (badha) of the empirically real occurs

when the absolute reality of the Brahman is realised. Thus, according to Vedanta, appearance implies the real, while the real need not imply appearance. To appear is essentially to appear in place of the real, but to be real is not necessarily to appear. All things -exist because the absolute exists. It is their Being. Thus the very existence of phenomena implies their non-existence as independent realities.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^\infty \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9 \left( U_9 - V_9 \int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9 \right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9 \int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$then given the constraints of $f(T_9), (A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 -$$

V 9M 9–]**e**–*T*9λ 9=0

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9 Z_9 \sin(Z_9 T_9) - (A_9 M_9 + Y_9) + (A_9 M_9 + X_9)\cos(Z_9 T_9) = 0$$

$$(A_9 M_9) Z_9 - M_9 Z_9 \cos(Z_9 T_9) + (A_9 M_9 - X_9)\sin(Z_9 T_9) = 0$$

Quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High

Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) **Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP** This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) **Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP** In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference: Journal of Physics: Conference Series 70 **Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows:** System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol

Where $A_9 = U_9 - V_9$Using (A2.3), substitute for $\boldsymbol{Z}_9^2$ and $Z_9$in (A2.2):

$W_9^2 (X_9^2 - 3 M_9 A_9 X_9 - M_9^2 (X_9 - Y_9) + 2 M_9^3 A_9 + 2 A_9^2 M_9^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_9$=1-cos ($Z_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (X_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ X_9 - Y_9/2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \leq W_9 \leq 2$ *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (X_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\ M_9)^2\ (2A_9 M_9 - (X_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9 M_9^2 - 3M_9 X_9\right) + W_9\left(3M_9 A_9^2 - 2A_9 X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9+Y_9\right)$$

$$=2M_9^2W_9^2(M_9+\mathrm{A}_9)+M_9A_9^2W_9+\left(2A_9M_9-X_9+Y_9\right)(A_9W_9+M_9W_9^2$$

$$+(2-W_9)((2M_9X_9+M_9Y_9)W_9+A_9(X_9+Y_9)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range 0 $W_9$ 2.Defining

$Q_9\left(W_9,A_9\right)=a_9W_9^2+c_9$ *where* $a_9b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

To Heidegger writings we shall compose the model. Following is not still written model able form. For one of the models please see any one of the papers under the appellation: Ïn Excited Reveries

In his later writings Heidegger drops the term 'phenomenon' in preference for the verbal form 'phainesthai' in order to emphasize even more the actuality or presentational property of Being. Explaining this new form of the term he writes: "Being disclosed itself to the ancient Greeks as 'physis'. The etymological roots 'phy-' and 'pha-' designate the same thing: 'phyein', the rising-up or upsurge which resides within itself as 'phainesthai', lighting-up, self-showing, coming-out, appearing-forth." Heidegger contrasted his notion of phenomenon with semblance (Schein) and with appearing (Erscheinung). In the case of semblance a thing can show itself as that which it is not, as when fool's gold shows itself to be gold. The ancients always allied semblance with non-being. Heidegger points out, however, that semblances are grounded in

showings, and so does Abhinava. Both Heidegger and Abhinava consequently maintain that all semblances have a real basis and are to be treated as instances of phenomena along with the so-called real showing or manifestation of non-deceptive objects. So Heidegger states that: 'how- ever much seeming, just that much being'. Thus self-showing or appearing defines Being as phenomenon, but this definition of Being is as yet incomplete. Being is not only self-showing but 'logos' which Heidegger explains means 'discourse' (Rede) in the sense of 'apophansis': 'letting-be-seen'. Phenomenology, which according to Heidegger is the only correct study of Being, means 'letting-be-seen-that-which- shows-itself. This is true of Saiva Paramadvaita as well. The reality of the world demands recognition; we are forced to accept the direct presentation of the fact of our daily experience. As Abhinava says: "if practical life, which is useful to all persons at all times, places and conditions were not real, then there would be nothing left which could be said to be real" A thousand proofs could not make 'blue' other than the colour blue. The reality of whatever appears in consciousness cannot be denied. Objects appear; they do not cease to do so by a mere emphatic denial. The manifestation of an entity in its own specific form is a fact at one level of consciousness; it is real. The appearing of the same entity in the same form but recognised to be a direct representation of the absolute is also a fact, but at another level of consciousness. It is no more or less real than the first. 'As is the state of consciousness, so is the experience,' says Abhinava. Although the nature of the absolute is discovered at a higher level of consciousness, The Doctrine of Vibration nonetheless it presents itself to us directly in the specific form in which we perceive things; otherwise there would be no way in which we could penetrate from the level of appearing to that of its source and basis. Abhinava writes: Real is the entity iyastu) that appears in the moment of direct perception (sak$atkara), that is to say, within our experience of it. Once its own specific form has been clearly determined one should, with effort, induce it to penetrate into its pure conscious nature. All things are known to be just as they present themselves. The concrete actuality of

being known (pramiti), irrespective of content, is itself the vibrant (spanda) actuality of the absolute. Liberating knowledge is gained not by going beyond appearances but by attending closely to them. "The secret," Mahesvarananda says, "is that liberation while alive (jivanmukti) is the profound contemplation of Maya's nature." No ontological distinction can be drawn between the absolute and its manifestations because both are an appearing (dbhasa), the latter of diversity and the former of 'the true light of consciousness which is beyond Maya and is the category Siva'. Those who have attained the category of Pure Knowledge above Maya and have thus gone beyond the category of Maya, see the entire universe as the light of consciousness . . . Just as the markings [on a feather] are nothing apart from the feather, the feather [is nothing apart from] them, similarly, when the light of consciousness is manifest, the whole group of phenomena is manifest as the light of consciousness itself. Within the sphere of Maya, every entity's 'own nature' (svabhdva) corresponds to its specific manifest form. Accordingly it is defined as that which distinguishes it from all else and from which it never deviates. Above the sphere of Maya, that is, above the level of objectivity, is the domain of the subject. At this level, everything is realised to be part of the fullness of the experience r and hence no longer bound by the conditions which impinge on the object. Here the part is discovered to be the whole, that is, consciousness in toto. In this sphere beyond relative distinctions, the yogi realises that (all) the categories of existence are present in every single category. m The yogi experiences every individual particular as the sum total of everything else. He recognises that all things have one nature and that every particular is all things. This is the 'essence' (sard) or co-extensive unity (samarasya) of all things. Integral Monism of Kashmiri SaivismWe have established that reality is manifest according to how [and the degree in which] the freedom of consciousness reveals it and that [this freedom] is the womb of all forms. Just as 'sweetness' is present in its entirety in every atom of the sugarcane, so each and every atom [of the universe] bears within itself the emanation of all things. This is

the level of consciousness in which the absolute reflects on itself realising to its eternal delight and astonishment (camatkara) its own integral nature. The reality of the world of diversity is not denied, but experienced in a new mode of awareness free of time and space in the eternal omnipresence of the Here and Now. [Phenomenal forms of awareness] such as 'this [exists]', born of the colouring [imparted to the absolute] by the limitations engendered by the diversifying power of time {kalakalana) also emanate within the Supreme Principle. There [at that level], Fullness {purriata) is the one nature [of all things] and so everything is omnipresent; otherwise, associated with division (khancjlana), the Fullness [of the absolute] would not be full. The content of absolute consciousness consists of diverse appearings (abhdsa) which, because they are manifest through it in this way, do not compromise the wholeness of consciousness. Everything we perceive is a momentary collocation of a number of such manifestations which combine together like 4 a row of altar lamps' (dipavalT) to form the single radiant picture of the universe. The individual objects which constitute the universe are specific collocations of such 'atomic' appearings. Together they form a single unified particular which appears according to its own defining features (svalaksand). A jar, for example, consists of a number of appearances such as 'round', 4 fat\ 'earthen', 'red', etc., which together discharge a single function (arthakriya), in this case, that of carrying the appearance 'water'. They unite with each other much as the scattered rays of a lamp come together when focused, or as the various currents of the sea together give rise to waves. Atomic appearings can combine in any number of ways, provided that they are not contrary to one another as established by the dictates of natural law (niyati). An appearance of 'form', for example, cannot combine with that of 'air'. Insofar as they share a common basis (samanyadhikaranya), a given cluster of appearances appears as a single whole. This common basis is the most prominent member of the group; the appearance 'jar' is such in the example quoted above. Any one appearance in a cluster may assume a more important or subordinate role.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_9 \square - b_9 \square 0$$

And

$$b_9^2 \square 4a_9 c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9 in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\ X_9$ ,

$Let\ Y_9 = \ y_9\ X_9$And try the substitution $A_9 = X_9/(2 + A_9\ y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + \ n_9 y_9)$ we have

$$W_9^2\ y_9\left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^2\left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9\,X_9\left[\frac{X_9\left(1+n_9y_9\right)}{M_9^2\left(2+n_9y_9\right)^2}+2\,y_9\left(n_9-1\right)+\frac{X_9^2\left(9+y_9\right)}{M_9^2\left(2+n_9y_9\right)}\right]=0.\ (A2.12)$$

Since $n_9$ and $A_9$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9=\left(2+n_9y_9\right),$ *the condition*$(A2.99?)$*can be written as*

$$\left[\frac{X_9\left(E_9-9\right)}{E_9^2}\right]^2\left[\frac{X_9^2}{M_9^2}-4\,y_9\,E_9^2\frac{X_9}{M_9^2}+4y_9\left(n_9-2\right)E_9^3\right]\ \square 0.$$

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9/W_9^2)$ can be expressed as

$$2\,E_9\left[y_9E_9+\left(y_9\,E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}\right]>\frac{X_9}{M_9^2}>2E_9\left[y_9E_9-y_9E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}]$$

dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique

for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) **Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP** Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) **Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP** Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a

cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) **Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP** Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service

placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) **Dynamic Service Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP** Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by $H*(X,(\wedge*TX)\otimes g)$, where X is the target hyperk\"ahler manifold. We further give a mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) **A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van**

**Raamsdonk1 and Nathan Seiberg2** study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. **Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP** Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy **The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP** Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the

present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. **ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP** Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition. Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb) the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537

© 2016 American Physical Society **Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP** Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant ($\Delta\alpha/\alpha$) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and $\Delta\alpha/\alpha$ to (e&eb) this conclusion. However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts $z\lesssim 3$, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the $\Lambda$ cold dark matter model value ($j\Lambda CDM=1$); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in $\alpha(z)$, linked to (e&eb) a linear coupling ($\zeta$) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order $|\zeta|\sim 10-4$, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in $\Delta\alpha/\alpha$ (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society **The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP** Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation

rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void

Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems, authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society **Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP** Authors measure the star formation properties of (e) two large samples of

galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the **star formation** rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a **function of environment**. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]**.** Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size $\sqrt{n} \times \sqrt{n}$? The problem is that $\sqrt{n}$ time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time $O(\sqrt{n})$ for $d \geq 3$, or $O(\sqrt{n}\log 5/2n)$ for $d = 2$. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search

tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ∗This work was mostly done while the author was a **PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship and by CIAR**. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 **Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP** Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can

lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 **Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP** Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. **Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP** Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e)

sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) **Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP** The relationship between the foundations of mathematics and (e&eb) physics is a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb) flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling

field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) **Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP** The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar

and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph] (or arXiv: 1503.05600v1 [math-ph] for this version) **Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP** This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs

mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) **Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP** Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and

derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to **the fine structure constant**. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) **Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP** Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor goes to 0 as the (e&eb) time goes to 0, the

big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph] (or arXiv: 1306.4613v1 [quant-ph] for this version) **Effects of gauge theory based number scaling on geometry Paul Benioff STOP** In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x,

outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph]

(or arXiv: 1211.3381v2 [quant-ph] for this version) **Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP** Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value,

\int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances, the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) **Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP** The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include

(e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) **Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP** Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) **New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP** This work is based on (e) the field of reference frames based on (e) quantum

representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects:Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc) Journal reference:Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) **A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul B**

**ABSOLUTE STABILITY**

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 ? $W_9$ ? 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9-9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2-2y_9\frac{(E_9+n_9^2y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right)-4y_9(E_9-n_9)$$ (A2.14)

And has the solution

$$\frac{X_9}{M_9^2}>\frac{E_9}{(E_9-9)}[y_9(E_9+n_9^2y_9)+[y_9^2(E_9+n_9^2y_9)^2+4y_9(E_9-9)(E_9-n_9)]^{9/2}]$$

Given $E_9>n_9$this lower limit to $X_9/M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9/M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9<n_9$ $(ie.,2A_9M_9^2<X_9-Y_9)$then (A2.9)and (A2.8)show that $Q_9$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2}>2E_9[y_9E_9-(y_9E_9(9+y_9)(E_9-n_9)]^{9/2}]$$

And

$$E_9>n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2/2M_9(X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$ and $X_9$ and $Y_9$ by $W_9^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

**Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability

of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e) long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007 STOP. David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer P** examine the possibility of coherent reversible information transfer between solid-state superconducting qubits and (e&eb) ensembles of ultracold atoms. Strong coupling between these systems is mediated by (e&eb) a microwave transmission line resonator that interacts near resonantly with (e&eb) the atoms via (e&eb) their optically excited Rydberg states. Solid-state qubits can then be used to (e) implement rapid quantum logic gates, while (e) collective metastable states of (e) the atoms can be employed for (e) long-term storage and (e&eb) optical readout of quantum information. **Reversible state transfer between superconducting qubits and (e&eb) atomic ensembles David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer Phys. Rev. A 79, 040304(R) – Published 14 April 2009 STOP.** Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures have been proposed. An

attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. **Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier** demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in (eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. **Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP.** In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown

to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , **Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally** frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. **Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg** propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states

into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. **Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011 http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan** proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 **Qubit transfer between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller** propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. **Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys.**

**Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich** report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. **Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi** investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. **Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev** investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer.

Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. **L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP.** In recent years **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P have** investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P** apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys. Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. **Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S.**

**Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP.** Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert **Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe** realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.**220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2** [quant-ph] for this version) **Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP.** Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. **Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker** describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue

from over (e&eb) a thousand plant and animal samples. **NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP.** Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this behavior, enhancing and accelerating (eb, eb+) the growth of correlations. **Michael Kastner** gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. **Michael Kastner** tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) **Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP.** A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. **Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E. Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91,**

**075435 – Published 26 February 2015 STOP.** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still

problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to

(e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP** Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that

(eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035

**Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime**

**C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013**

**STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and

cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP** Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex

Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) **A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen**

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9 and$ separate (A3.1)into real and imaginary parts

$$Z_9^2 + \left(M_9 + V_9\right) Z_9 \sin\left(Z_9 T_9\right) - \left(U_9 M_9 + Y_9\right) + (M_9\left(U_9 + V_9\right) - X_9)$$

$$\mathrm{X}\cos\left(Z_9 T_9\right) - 2M_9 V_9 cos^2\left(Z_9 T_9\right) + M_9 V_9 = 0$$

$$(M_9 + U_9)\, Z_9 - (M_9 + V_9)\, Z_9\ cos\left(Z_9 T_9\right) + \left(M_9\left(U_9 + V_9\right) - X\right)$$

$$\mathrm{X}\sin\left(\left(Z_9 T_9\right) - M_9 V_9 \sin\left(Z_9 T_9\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 =$ $1 - cos\, Z_9\ T_9$and $A_9 = U_9 - V_9$we obtain

$$4\, M_9^2 V_9^2 W_9^4 + 2\, M_9 V_9\ W_9^3[(M_9 - V_9)^2 + A_9\left(M_9 + V_9\right) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2\left[M_9 - V_9\right)^2\left(2\, A_9\ M_9 + Y_9 - X_9\right) + \left(2A_9\ M_9 - X_9\right)^2 - A_9\left(M_9 - V_9\right)\left(2A_9\ M_9 - X_9\right)$$

$$+ M_9 V_9\,(A_9^2 + 4\left(X_9 + Y_9\right))]$$

$$+ W_9\left[A_9^2 + 2X_9\right)\left(2\, A_9\ M_9 - X_9\right) + 2A_9 Y_9\,(M_9 + V_9) - M_9 A_9^3] + A_9^2\left(X_9\ M_9 - Y_9\right) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9\left(2, U_9 V_9\right) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting cos $(Z_9 T_9) = -1\ in\ (A3.3)$

We have applied the model to Kashmiri Shaivism. Whose elaborate rituals is enough to make a indivual raise from individual consciousness to cosmic consciousness which we have called Cosmic Consciousness.

It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) **Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP** The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic

Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** The group of unimodular transformations on the roots of polynomials over GF(2) is considered, and those polynomials with (e&eb) symmetries in (eb) the unimodular group are identified. The cross-correlation function between (e&eb) two maximum-length linear shift register sequences of (e) the same degree is shown to be (=) computable as an explicit linear transformation, in (eb) matrix form, on (e&eb) either one of the sequences, regarded as (=) a vector. The underlying vector space is (=) the "cyclotomic algebra," generated by (e) the cyclotomic cosets, or by what Gauss termed the **"periods" of the cyclotomic equation**. A variety of numerical examples are worked in detail. **Information Sciences Volume 1, Issue 1, December 1968, Pages 87-109 Theory of transformation groups of polynomials over GF(2) with applications to linear shift register sequences Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more link. Solomon W. Golomb** stop Authors study the P−V criticality and phase transition in (eb) the extended phase space of charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of (e) the system and its conjugate quantity is (=) the thermodynamic volume of the black hole. The black holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb) no P−V criticality and phase transition appear, while (e) for the black holes with a spherical horizon, even when the charge of the black hole is absent, t(eb) he P−V criticality and the small black hole/large black hole phase transition will appear, but it happens only in d=5 dimensions; when the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to have (e) the P−V criticality

and (e&eb)the small black hole/large black hole phase transition, there exists (ebb) an upper bound for the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the **Gauss-Bonnet coefficient** and Q is (=) the charge of the black hole. Authors calculate the critical exponents at the critical point and find that (eb) for all cases, (eb) they are the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09(2013)005 Cite as: arXiv:1306.6233 [gr-qc] (or arXiv:1306.6233v4 [gr-qc] for this version) **P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in IR but with (e&eb) improved UV behaviors. Here, authors study an IR modification which breaks ``softly'' (e) the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Authors obtain the black hole and cosmological solutions for **``arbitrary'' cosmological constant** that represent (eb) the analogs of the standard Schwarzschild-(A)dS solutions which can be asymptotically (A)dS as well as flat and authors discuss their thermodynamical properties. They also obtain solutions for (e) FRW metric with (e&eb) an arbitrary cosmological constant. Authors study its implication to (e&eb) the dark energy and find that it seems to be consistent with (e&eb) current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP.** The Einstein equations with (e&eb) a negative cosmological constant admit (e) black hole solutions which are asymptotic to (e) anti-de Sitter space. Like black holes in asymptotically flat space, these solutions have (e) thermodynamic properties including (e) a characteristic temperature and intrinsic entropy equal to (=) one quarter of the area of the event horizon in Planck units. There are however some important differences from (e) the asymptotically flat case. A black hole in anti-de Sitter space has (e) a minimum temperature which occurs when (e) its size is of the order of the characteristic radius of (e) the anti-de Sitter space. For larger black holes the red-shifted

temperature measured at (eb) infinity is greater. This means that such black holes have (e) positive specific heat and can be in stable equilibrium with (e&eb) thermal radiation at a fixed temperature. It also implies (eb) that the canonical ensemble exists for asymptotically anti-de Sitter space, unlike the case for asymptotically flat space. One can also consider the microcanonical ensemble. One can avoid (e) the problem that arises in asymptotically flat space of (e) having to put the system in (eb) a box with unphysical perfectly reflecting walls because (e) the gravitational potential of anti-de Sitter space acts as a box of finite volume. Communicated by A. Jaffe **Communications in Mathematical Physics December 1983, Volume 87, Issue 4, pp 577-588 Thermodynamics of black holes in anti-de Sitter space S. W. Hawking, Don N. Page STOP** In previous work authors introduced the concept of black-hole entropy, which they identified with (e&eb) the surface area of the black hole in question expressed in units of (e) the Planck length squared. Authors suggested that the appropriate generalization of the second law for a region containing a black hole is (=) that the black-hole entropy plus (+) the common entropy in the black-hole exterior never decreases. Here authors establish the validity of this law for the infall of an entropy-bearing system into (e&eb) a much larger and more massive generic stationary black hole. To do this authors determine (eb) a general lower bound for the increase in black-hole entropy, and an upper bound for the entropy of the system, while allowing for quantum effects at each stage. In passing authors show that the generalized second law is a statistical law which becomes (=) over-whelmingly probable in the limit of a macroscopic system. Authors also consider briefly more general situations. Finally, we give two simple examples of predictions made by the generalized second law for (e) black-hole formation processes. Received 17 September 1973 DOI:http://dx.doi.org/10.1103/PhysRevD.9.3292 **Generalized second law of thermodynamics in black-hole physics Jacob D. Bekenstein Phys. Rev. D 9, 3292 – Published 15 June 1974 STOP** The Lovelock gravity consists of (e) the dimensionally extended Euler densities. The geometry and horizon structure of (e) black hole solutions could be quite complicated in (eb) this gravity; however, authors find that some thermodynamic quantities

of the black holes like the mass, Hawking temperature and entropy have (e) simple forms expressed in terms of (e) horizon radius. The case with black hole horizon being a Ricci flat hypersurface is (=) particularly simple. In that case the black holes are (=) always thermodynamically stable with (e&eb) a positive heat capacity and their entropy still obeys (e&eb) the area formula, which is no longer valid for (e) black holes with positive or negative constant curvature horizon hypersurface. In addition, for black holes in (eb) the gravity theory of Ricci scalar plus a 2n-dimensional Euler density with (e&eb) a positive coefficient, thermodynamically stable small black holes always exist in (eb) D=2n+1 dimensions, which are absent in (eb) the case without (e) the Euler density term, while (e) the thermodynamic properties of the black hole solutions with (e&eb) the Euler density term are qualitatively similar to those of black holes without (e) the Euler density term as D>2n+1. Copyright © 2004 Published by Elsevier B.V. **Physics Letters B Volume 582, Issues 3–4, 4 March 2004, Pages 237–242 A note on thermodynamics of black holes in Lovelock gravity Rong-Gen CAI STOP** There are a number of similarities between black-hole physics and (e, eb, =) thermodynamics. Most striking is the similarity in the behaviors of black-hole area and (e&eb) of entropy: Both quantities tend to increase (eb+) irreversibly. In this paper we make this similarity the basis of (e) a thermodynamic approach to black-hole physics. After a brief review of the elements of the theory of information, authors discuss black-hole physics from the point of view of information theory. Authors show that it is natural to introduce (eb) the concept of black-hole entropy as (=) the measure of information about (e&eb) a black-hole interior which is (=) inaccessible to an exterior observer. Considerations of simplicity and consistency, and dimensional arguments indicate (eb) that the black-hole entropy is equal to (=) the ratio of the black-hole area to the square of the Planck length times a dimensionless constant of (e) order unity. A different approach making use of the specific properties of Kerr black holes and of concepts from information theory leads to (eb) the same conclusion, and suggests (eb) a definite value for the constant. The physical content of (e) the concept of black-hole entropy derives from (e) the following generalized version of the second law: When common entropy goes down a

black hole, (eb) the common entropy in (eb) the black-hole exterior plus the black-hole entropy never decreases. The validity of this version of the second law is supported by (e) an argument from information theory as well as by several examples. Received 2 November 1972 DOI:http://dx.doi.org/10.1103/PhysRevD.7.2333 **Black Holes and Entropy Jacob D. Bekenstein Phys. Rev. D 7, 2333 – Published 15 April 1973 STOP** The effects of power-law Maxwell field on the Van der Waals like phase transition of higher-dimensional dilaton black holes are probed in detail. It is shown that the Smarr relation gains corrections due to (e) the effects of (e) both the power-law Maxwell field and (e&eb) the dilaton field while (e) thermodynamic volume is exactly the same as (=) that of Einstein-Maxwell-dilaton black holes. Authors successfully derive the analytic solutions of (e) critical point and carry out (e&eb) some check to ensure that these critical quantities are (=) positive. It is shown that the constraint on the parameters turns out to be (=) 0<\alpha^2<1, which is tighter than that in (eb) the non-extended phase space. It is also shown (eb) that these critical quantities and (e&eb) the ratio P_cv_c/T_c are affected by (e&eb) the power-law Maxwell field. Moreover, critical exponents are found to coincide with those of other AdS black holes, showing (eb) the powerful influence of mean field theory. Comments: references added Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1601.05500 [gr-qc] (or arXiv:1601.05500v2 [gr-qc] for this version **Effects of power-law Maxwell field on the Van der Waals like phase transition of higher dimensional dilaton black holes Jie-Xiong Mo, Gu-Qiang Li, Xiao-Bao Xu STOP** A consistent and unified picture for (e) critical phenomena of charged AdS black holes in f(R) gravity is drawn in this paper. Firstly, we investigate the phase transition in canonical ensemble. Authors derive the explicit solutions corresponding to (e&eb) the divergence of CQ. The two solutions merge into (e&eb) one when (e) the condition Qc=−13R0−−−√ is satisfied. The curve of specific heat for Q<Qc has two divergent points and can be divided into (e&eb) three regions. Both the large radius region and the small radius region are (=) thermodynamically stable with (e&eb) positive specific heat while (e) the medium radius region is unstable with (e&eb) negative specific heat. However, when Q>Qc, the specific heat is

always positive, implying (eb) the black holes are locally stable and no phase transition will take place. Secondly, both the T−r+ curve and T−S curve f(R) AdS black holes are investigated and they exhibit Van der Vaals like behavior as the P−v curve in the former research. Critical physical quantities are obtained and they are consistent with (=) those derived from the specific heat analysis. Authors carry out numerical check of Maxwell equal area law for the cases Q=0.2Qc, 0.4Qc, 0.6Qc, 0.8Qc. The relative errors are (=) amazingly small and can be negligible. So the Maxwell equal area law holds for (e) T−S curve of f(R) black holes. Thirdly, authors establish geometrothermodynamics for (e) f(R) AdS black hole to examine the phase structure. It is shown that the Legendre invariant scalar curvature R would diverge exactly where (e) the specific heat diverges. To summarize, the above three perspectives are consistent with (e&eb) each other, thus providing (eb) a unified picture which deepens the understanding of (e) critical phenomena of (e) f(R) AdS black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1602.01251 [gr-qc] (or arXiv:1602.01251v1 [gr-qc] for this version) **A consistent and unified picture for critical phenomena of f(R) AdS black holes Jie-Xiong Mo, Gu-Qiang Li, Yu-Cheng Wu STOP** The state equation of the charged AdS black hole is reviewed in (e&eb) the T−r plane. Thinking of the phase transition, the T−S, P−V, P−v graphs are plotted and then the equal area law is used (e) in the three cases to get (eb) the phase transition point (P, T). The analytical phase transition point relations for (e) P-T of charged AdS black hole has been obtained successfully. By comparing the three results, authors find that the equal area law possibly cannot (e) be used directly for P−v plane. According to the T−S, P−V results, they plot (e&eb) the P−T−Q graph and find that for a highly charged black hole a very low temperature condition is required for (e) the phase transition. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI:10.1140/epjc/s10052-015-3641-0 Cite as: arXiv: 1503.07658 [gr-qc] (or arXiv:1503.07658v1 [gr-qc] for this version) **A note on Maxwell's equal area law for black hole phase transitionShan-Quan Lan, Jie-Xiong Mo, Wen-Biao Liu STOP** Recently, extended phase space thermodynamics of Lovelock AdS black holes has been of great interest. To provide

insight from (e) a different perspective and gain (e) a unified phase transition picture, non-extended phase space thermodynamics of (e) (n+1)-dimensional charged topological Lovelock AdS black holes is (=) investigated detailedly in (eb) the grand canonical ensemble. Specifically, the specific heat at (eb) constant electric potential is calculated and phase transition in the grand canonical ensemble is discussed. To probe the impact of the various parameters, they utilize (e) the control variate method and solve (e) the phase transition condition equation numerically for the case k=1, −1. There are two critical points for (e) the case n=6, k=1 while there is only one for other cases. For k=0, there exists no phase transition point. To figure out the nature of phase transition in the grand canonical ensemble, authors carry out an analytic check of the analog form of (e) Ehrenfest equations proposed by Banerjee et al. It is shown that Lovelock AdS black holes in (eb) the grand canonical ensemble undergo (e&eb) a second order phase transition. To examine the phase structure in (eb) the grand canonical ensemble, authors utilize the thermodynamic geometry method and calculate both the Weinhold metric and Ruppeiner metric. It is shown that for both analytic and graphical results that the divergence structure of (e) the Ruppeiner scalar curvature coincides with (=) that of the specific heat. Research provides (eb) one more example that Ruppeiner metric serves as (=) a wonderful tool to (e) probe the phase structures of (e) black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C75 (2015) 5, 211 DOI:10.1140/epjc/s10052-015-3434-5 Cite as: arXiv: 1503.01956 [gr-qc] (or arXiv: 1503.01956v1 [gr-qc] for this version **Non-extended phase space thermodynamics of Lovelock AdS black holes in grand canonical ensemble Jie-Xiong Mo, Wen-Biao Liu STOP** To provide an analytic verification of the nature of phase transition at the critical point of P−V criticality, the original expressions of Ehrenfest equations have been introduced directly. By treating the cosmological constant and (e&eb) its conjugate quantity as (=) thermodynamic pressure and volume respectively, authors carry out analytical check of (e) classical Ehrenfest equations. To show that approach is universal, authors investigate not only higher-dimensional charged AdS black holes (e&eb), but also rotating AdS black holes. Not only are the examples

of Einstein gravity shown, but also the example of modified gravity is presented for (e) Gauss-Bonnet AdS black holes. The specific heat at constant pressure CP, the volume expansion coefficient α and the isothermal compressibility coefficient κT are found to diverge exactly at the critical point. It has been verified that both Ehrenfest equations hold at (eb) the critical point of P−V criticality in (eb) the extended phase spaces of (e) AdS black holes. So the nature of the critical point of P−V criticality of AdS black holes has been demonstrated analytically to be (=) a second-order phase transition. These results are consistent with (e&eb) the nature of liquid-gas phase transition at the critical point. In this sense, extant research would deepen the understanding of (e) the relations of AdS black holes and (e&eb) liquid-gas systems. Moreover, successful approaches to introduce (eb) the original expressions of Erhenfest equations directly into (e&eb) black hole phase transition research demonstrate again (eb) that black hole thermodynamics is closely related (e&eb) to classical thermodynamics, which allows (eb) us to borrow techniques from (e) classical thermodynamics to investigate the thermodynamics of black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 89, 084057 (2014) DOI: 10.1103/PhysRevD.89.084057 Cite as:arXiv: 1404.3872 [gr-qc] (or arXiv:1404.3872v1 [gr-qc] for this version) **Ehrenfest scheme for P−V criticality of higher dimensional charged black holes, rotating black holes and Gauss-Bonnet AdS black holes Jie-Xiong Mo, Wen-Biao Liu STOP** Aiming at a unified phase transition picture of (e) the charged topological black hole in (eb) Ho\v{r}ava-Lifshitz gravity, authors investigate this issue not only in (eb) canonical ensemble with (e&eb) the fixed charge case but also in grand-canonical ensemble with (e&eb) the fixed potential case. Authors firstly perform the standard analysis of the specific heat,(e&eb) the free energy and the Gibbs potential, and (e&eb) then study its geometrothermodynamics. It is shown that the local phase transition points (eb) not only witness the divergence of (e) the specific heat, but also witness the minimum temperature and the maximum free energy or Gibbs potential. They also witness the divergence of (e) the corresponding thermodynamic scalar curvature. No matter which ensemble is chosen, the metric constructed can

successfully produce (eb) the behavior of the thermodynamic interaction and (e&eb) phase transition structure while other metrics failed to (e) predict the phase transition point of (e) the charged topological black hole in former literature. In grand-canonical ensemble, authors have discovered the phase transition which has not been reported before. It is similar to (e) the canonical ensemble in which the phase transition only takes place when k=−1. But it also has its unique characteristics that the location of the phase transition point depends on (eb) the value of potential, which is different from (e) the canonical ensemble where (e) the phase transition point is (=) independent of the parameters. After an analytical check of Ehrenfest scheme, authors find that the new phase transition is (=) a second order one. It is also found that the thermodynamics of (e) the black hole in Horava-Lifshitz gravity is quite different from that in (eb) Einstein gravity. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP1310(2013)056 DOI:10.1007/JHEP10(2013)056 Cite as: arXiv:1404.2497 [gr-qc] (or arXiv:1404.2497v1 [gr-qc] for this version) **A unified phase transition picture of the charged topological black hole in Horava-Lifshitz gravity Jie-Xiong Mo, Xiao-Xiong Zeng, Gu-Qiang Li, Xin Jiang, Wen-Biao Liu STOP** To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case $\beta\to\infty$ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness condition of (e) entropy. It is shown

that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure authors have discovered. It is shown that Legendre invariant thermodynamic scalar

curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws, leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions, geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged AdS black holes by considering the case of the RN-AdS black holes

surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase transition. For the case $\omega_q=-2/3$ authors derive analytic expressions of (e) critical physical quantities, while (e) for cases $\omega_q\neq-2/3$ appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this

version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=) chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy

Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed

quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions, while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma?**

All this consists of the act of consciousness alone", says Vasubandhu, "because unreal entities appear, just as a man with defective vision sees unreal hair or a moon, etc."He points to dreams as examples of purely subjective constructs which appear to be objective realities. The apparent reality dreams possess is not derived from any concrete, objective world, but merely from the idea of objectivity. While the Yogacara does not say that an idea has, for example, spatial attributes, it does have a form manifesting them. While he agrees with the Saiva idealist that

appearances have no independent existence apart from their appearing to consciousness, he maintains that for this reason they are unreal. The creativity of consciousness consists in its diversification in many modes having apparent externality; it is not a creation of objects. While the Kashmiri Saivite agrees that the world is pure conscious- ness alone, he maintains that it is such because it is a real creation of consciousness. The effect is essentially identical with the cause and shares in its reality. Matter and the entire universe are absolutely real, as 'congealed' (sty ana) or 'contracted' (samkucita) forms of consciousness. "This God of consciousness", writes Ksemaraja, "generates the universe and its form is a condensation of His own essence (rasa)" m By boiling sugarcane juice it condenses to form treacle, brown sugar and candy which retains its sweetness. Similarly, consciousness abides unchanged even though it assumes the concrete material form of the five gross elements. The same reality thus abides equally in gross and subtle forms. Consequently no object is totally insentient. Even stones bear a trace (vasana) of consciousness, although it is not clearly apparent because it is not associated with the vital breath (prana) and other components of a psycho-physical organism. Somananda goes so far as to affirm that physical objects, far from being insentient, can only exist insofar as they are aware of themselves as existing. The jar performs its function because it knows itself to be its agent. Indeed, all things are pervaded by consciousness and at one with it and hence share in its omniscience. Thus, Siva, Who perceives Himself in the form of physical objects, is the one ultimate reality. "The jar knows because it is of my nature", writes Somananda, "and I know it because I am of the jar's nature. I know because I am of Sadasiva's nature and He knows because He is of my nature; Yajnadatta [knows] because he is of Siva's nature and Siva [knows] because He is of Yajnadatta's nature". Integral Monism of Kashmiri Saivism Everything in this sense is directly perceived by absolute conscious- ness, and this direct perception (pratyaksa) unifies the knowable into a single, undivided whole. This is the central concept behind a doctrine originally expounded by Narasimha called 'the non-dualism of direct perception' (pratyaksadvaita). This states that consciousness is essentially perceptive and that its perception of all things operates throughout the universe. Insofar

as phenomena are clearly evident (sphufa) to us, everything is directly perceived by absolute consciousness, with which our individual consciousness is identical. This direct perception unfolds everywhere; the one true reality, it is alone and without companion or rival (nihsapatna). Even though it remains one, it can, by its very nature, perceive distinctions (bheda) between one entity and another, without this engendering any division within it. We distinguish between two entities in empirical terms on the basis of their mutual exclusion (anyonyabhava). The relative distinction {bheda) between them is essentially the perceived difference between their respective characteristics. Despite this difference they are united within the purview of a single cognition insofar as they are equally both manifest appearances. This cognition is the undivided essence (rasa) or 'own nature' (svabhava) of both. Encompassed by the 'fire of consciousness', there is no essential difference between them. Just as when an emerald and ruby reflect each other's light, the ruby is reddish-green and the emerald greenish-red, similarly everything is connected with everything else as part of the single variegated (vicitra) cognition of absolute consciousness. Mahesvarananda writes: The Supreme Lord's unique state of emotivity (asadharanabhava) is the outpouring of pure Being (mahasatta). It is manifest as the brilliance (sphuratta) of the universe which, if we ponder deeply, [is realized to be] the single flavour (ekarasa) of the essence of Beauty which is the vibration of the bliss of one's own nature. In this way all things are in reality one although divided from the one another sharing as they do the 'single flavour' (ekarasa) of the pure vibration of consciousness. Kashmiri Saiva Realism Kashmiri Saivism as a whole has been variously called a form of 'realistic idealism', 'monistic idealism', ‘idealistic monism' and 'concrete monism'. It is easy to understand why Kashmiri Saivism is The Doctrine of Vibration said to be 'idealistic' and 'monistic', but in what sense is it also 'realistic*? The answer to this question is of no small importance in trying to under- stand the central idea behind its metaphysics and the fundamental importance of the concept of Spanda, in this seemingly impossible marriage between monistic idealism and pluralistic realism. The Kashmiri Saiva approach understands the world to be a symbol of the absolute, that is, as the manner in which it presents itself to us. Again we can

contrast this view with that of the Advaita Vedanta. The Advaita Vedanta understands the world to be an expression of the absolute insofar as it exists by virtue of the absolute's Being. Being is understood to be the real unity which underlies empirically manifest separateness and as such is never empirically manifest. It is only transcendentally actual as 'being-in-itself. The Kashmiri Saiva position represents, in a sense, a reversal of this point of view. The nature of the absolute, and also that of Being, is onceived as an eternal becoming (satatodita), a dynamic flux or Spanda, 'the agency of the act of being'. It is dentified with the concrete actuality of the fact of appearing, not passive unmanifest Being. Appearance (abhasa) alone is real.ppearing (prakasamanatva) is equivalent to the fact of being (astitva). Ksemaraja writes in his commentary on the Stanzas on Vibration: Indeed, all things are manifest because they are nothing but manifestation. The point being that nothing is manifest apart from manifestation The absolutely unmanifest, from this point of view, can have as little existence as the space in a lattice window of a sky-palace. Nay, even less, because even that space can appear as an imagined mage manifest within consciousness. Everything is real according to the manner in which it appears. Even an illusion is in this sense real, insofar as it appears and is known in the manner in which it appears. The empirical and the real are identical categories of thought. As Abhinava says: Thus this is the supreme doctrine (upanisatf), namely that, when- ever and in whatever form [an entity] appears, that then is its particular nature. Perhaps at this stage a brief comparison with Heidegger's ideas might prove to be enlightening and not altogether out of place. According to Heidegger's phenomenology of Being, reality is intelligible in a two-fold Integral Monism of Kashmiri Saivism manner as 'phenomenon' and 'logos'. Heidegger defines what he means by 'phenomenon' as: "that-which-shows-itself. The manifest . . . phenomena are then the collection of that which lies open in broad daylight or can be brought to the light of day — what the Greeks at times implicitly identified as 'ta onta' (the things-which-are)".

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb) heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna**

**Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version) **Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb) strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel

period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar (=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI: 10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills (SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear

differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black**

**holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-

relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb)

the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this

version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstate thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11 August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of

entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.** Characterizing the behaviour of strongly coupled

quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at λ=λc, the correlation length diverges as ξ∼|λ−λc|−ν, while (e) the gap between the ground state and first excited state vanishes as (e) ΔE∼ξ−z∼|λ−λc|zν. Exactly at the critical point λ=λc, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as ω∼|k−k0|z. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-**

**equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004 STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element.

**Model has been applied for the following results:**

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| **Searches for Sterile Neutrinos with the IceCube Detector M. G. Aartsen et al. (IceCube Collaboration) Phys.**<br><br>**Rev. Lett 117, 071801 – Published 8 August 2016 Physics logo See Viewpoint: Hunting the Sterile Neutrino** |
| **(1) The IceCube neutrino telescope at the South Pole has measured the atmospheric muon neutrino spectrum as (=) a function of zenith angle and (e&eb) energy in the approximate 320 GeV to 20 TeV range, to search for (e) the oscillation signatures of (e) light sterile neutrinos.** |
| **(2) No evidence for anomalous νμ or ν¯μ disappearance is observed in (eb) either of two independently developed analyses, each using one year of (e) atmospheric neutrino data.** |
| **(3) New exclusion limits are placed on (e&eb) the parameter space of (e) the 3+1 model, in which muon antineutrinos experience (e) a strong Mikheyev-Smirnov-Wolfenstein-resonant oscillation.** |
| **(4) The exclusion limits extend to (e) sin22024≤0.02 at Δm2∼0.3 eV2 at the 90% confidence level.** |
| **(5) The allowed region from (e) global analysis of appearance experiments, including (e) LSND and MiniBooNE, is excluded at approximately the 99% confidence level for (e) the global best-fit value of \|Ue4\|2. Received 5 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.117.071801 © 2016 American Physical Society** |
| **Synopsis: Searching for Majorana Neutrinos August 16, 2016The KamLAND-Zen collaboration has run the most sensitive search to date for a radioactive decay that could reveal whether neutrinos are Majorana fermions. Synopsis figureKamLAND-Zen Collaboration** |
| **(6) Are neutrinos and (=) antineutrinos the same particles, as hypothesized by (e) physicist Ettore Majorana?** |
| **(7) If true, this "Majorana" nature of neutrinos could explain (eb) why neutrinos have (e) mass and why the Universe has (e) more matter than antimatter.** |
| **(8) Using an underground neutrino detector in Japan, the KamLAND-Zen collaboration has carried out the most sensitive search to date for a process called neutrinoless double-beta decay, which could confirm (eb) Majorana's hypothesis.** |
| **(9) The results set stringent constraints on the possible properties of (e) Majorana neutrinos.** |
| **(10) Double-beta decay is (=) a nuclear process in which two neutrons turn into (e&eb) two protons (or vice versa), emitting (eb) two antineutrinos.** |
| **(11) If neutrinos are Majorana fermions, (eb) the two antineutrinos could annihilate each other, and some double-beta decays would emit (eb) no antineutrinos.** |
| **(12) However, neutrinoless decay is believed to be (=) extremely rare. Searches of the process using (e) large amounts of the 35 naturally occurring (eb) isotopes that can undergo (e&eb) double-beta decay (such as xenon-136) have so far come up empty.** |

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| **(13) The KamLAND-Zen collaboration boosted their experimental sensitivity to (e) neutrinoless double-beta decays through (e&eb) a combination of factors.** |
| **(14) The researchers deployed clean detectors with (e&eb) very low background noise.** |
| **(15) And they used an unprecedented amount of xenon-136, which they purified to (e) remove any radioactive contaminants that could produce (eb) unwanted signals in the detectors.** |
| **(16) Combining data from several years, they improved (eb) the limit on the probability of neutrinoless double-beta decay by (e) sixfold compared to previous searches.** |
| **(17) This probability corresponds to (e&eb) a decay lifetime whose staggering value exceeds (e) 1026 years.** |
| **(18) This measurement allowed (eb) the researchers to set (eb) the most stringent upper limit on (e&eb) the possible mass of Majorana neutrinos (the particle must be lighter than 61–165 meV). This research is published in Physical Review Letters. –Matteo Rini Matteo Rini is the Deputy Editor of Physics.** |
| **How to Reach a Thousand-Second in-Plane Polarization Lifetime with 0.97−GeV/c Deuterons in a Storage Ring G. Guidoboni et al. (JEDI Collaboration) Phys. Rev. Lett. 117, 054801 – Published 28 July 2016 Physics logo See Synopsis: Deuterons Spin together for a Thousand Seconds** |
| **(19) Authors observe a deuteron beam polarization lifetime near 1000 s in (eb) the horizontal plane of (e) a magnetic storage ring (COSY).** |
| **(20) This long spin coherence time is maintained through (e&eb) a combination of beam bunching, electron cooling, (e&eb) sextupole field corrections, and (e&eb) the suppression of collective effects through (e&eb) beam current limits.** |
| **(21) This record lifetime is required for (e) a storage ring search for (e) an intrinsic electric dipole moment on (e&eb) the deuteron at a statistical sensitivity level approaching 10−29 e cm.** |
| **Synopsis: Deuterons Spin together for a Thousand Seconds July 28, 2016 Researchers set a new record for the in-plane spin-alignment lifetime of deuterons circulating in a magnetic storage ring. Synopsis figure E. Stephenson/Indiana University Center for Spacetime Symmetries Sime sentences and figures are deleted due to spatial restrictions.** |
| **(22) The electric dipole moment (EDM) of (e) a fundamental particle or a simple nucleus is believed to be (=) extremely small.** |
| **(23) One might therefore think that it isn't worth measuring this quantity. But it is. A nonzero EDM would signal (eb) a broken symmetry of (e) nature known as CP violation, which is required to (e) explain why the Universe has (e) more matter than antimatter.** |
| **(24) Edward Stephenson from Indiana University, Bloomington, and colleagues has now made a significant stride towards determining (eb) the EDM of the deuteron using (e) a deuteron beam circulating in (e&eb) a magnetic storage ring.** |

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| **(25) They have been able to preserve the alignment of the deuterons' spins for (e) a chunk of time that may be long enough to determine (eb) the deuteron's EDM.** |
| **(26) In experiments of this kind, the EDM can be determined by (e) measuring the rotation of (e) the initially aligned spins of (e) the beam particles as they travel around (e&eb) the ring.** |
| **(27) To obtain a measurement sufficiently accurate to explore (e&eb) CP violation predicted by (e) theories beyond the standard model, the spin alignment has to be (=) preserved for about 1000 s.** |
| **(28) But changes in particle speed—associated with (e&eb) position or direction variations among (e&eb) the particles in the beam—lead to (eb) a rapid loss of the alignment in the plane of the ring.** |
| **(29) Stephenson and co-workers have managed to retain (eb)the alignment of deuterons for the required 1000 s, which is 3 orders of magnitude longer (eb) than the previous record (obtained for an electron beam).** |
| **(30) Their success resulted from (e) capturing the beam into a tight bunch to reduce (e) position and direction spread and from (e) shaping the ring's magnetic field to(e) ensure all of the deuteron orbits around (e&eb) the ring have the same length, despite (e) direction variations. This research is published in Physical Review Letters. –Ana Lopes Ana Lopes is a Senior Editor of Physics.** |
| **Cotunneling Drag Effect in Coulomb-Coupled Quantum Dots A. J. Keller, J. S. Lim, David Sánchez, Rosa López, S. Amasha, J. A. Katine, Hadas Shtrikman, and D. Goldhaber-Gordon Phys. Rev. Lett. 117, 066602 – Published 4 August 2016 Physics logo See Synopsis: Coulomb Drag in a Double Dot** |
| **(31) In Coulomb drag, a current flowing in one conductor can induce (eb) a voltage across an adjacent conductor via (e&eb) the Coulomb interaction.** |
| **(32) The mechanisms yielding drag effects are not always understood, even though drag effects are (=) sufficiently general to be (=) seen in many low-dimensional systems.** |
| **(33) In this Letter, authors observe Coulomb drag in (eb) a Coulomb-coupled double quantum dot and, through (e&eb) both experimental and theoretical arguments, identify cotunneling as (=) essential to obtaining a correct qualitative understanding of the drag behavior.** |
| **Synopsis: Coulomb Drag in a Double Dot August 4, 2016 Electric current passing through a quantum dot can generate current in a nearby dot through a coordinated tunneling mechanism. Synopsis figureA. J. Keller et al., Phys. Rev. Lett. (2016)** |
| **(34) Electrons moving through (e&eb) a conductor can induce (eb) the movement of electrons of a nearby, but isolated, conductor.** |
| **(35) This so-called Coulomb drag is well understood in 2D layered materials, but the theory for nanowires and quantum dots does not (e) fully explain experiments.** |
| **(36) A new study of Coulomb drag in two nearby quantum dots shows (eb) that simultaneous tunneling of electrons—called cotunneling—can help explain the low-temperature behavior of (e) the system.** |

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| **(37) Electric current flows through (e&eb) a quantum dot by the tunneling of (e) electrons.** |
| **(38) When two dots are placed next to each other, the active, or lead, dot can drag current through (e&eb) the passive dot by (e) a process called sequential tunneling.** |
| **(39) In this case, an electron first tunnels onto (e&eb) the active dot, increasing (eb) the probability of tunneling onto (e&eb) the passive dot.** |
| **(40) Thus the passive dot follows (e&eb) the active dot's lead, but always one step behind.** |
| **(41) Sequential tunneling explains (eb) Coulomb drag observations at (eb) high temperature, but no studies have investigated cold temperatures where (e) quantum coherence effects may play a role.** |
| **(42) Andrew Keller from the California Institute of Technology, Pasadena, and his colleagues explored Coulomb drag at (eb) sub-Kelvin temperatures in a pair of quantum dots coupled by (e) capacitance.** |
| **(43) The team applied a current-driving voltage across (e&eb) one dot and detected a drag current in (eb) the other dot, even in regimes where (e) the driving voltage was very small.** |
| **(44) This conflicts with (e&eb) the sequential tunneling model, which predicts (eb) a threshold driving voltage below which (e) the drag current should be zero.** |
| **(45) The team showed that the inclusion of cotunneling into (eb) the drag theory explains the observations.** |
| **(46) Cotunneling could be important to efforts aimed at harvesting a unidirectional current from (e) a fluctuating voltage supply. This research is published in Physical Review Letters. –Michael Schirber Michael Schirber is a Corresponding Editor for Physics based in Lyon, France.** |
| **Maybe it wasn't the Higgs particle after all November 7, 2014 Maybe it wasn't the Higgs particle after all Associate Professor Mads Toudal Frandsen, University of Southern Denmark Credit: University of Southern Denmark** |
| **(47) Last year CERN announced the finding of a new elementary particle, the Higgs particle. But maybe it wasn't the Higgs particle, maybe it just looks like it. And maybe it is not alone. Many calculations indicate that the particle discovered last year in the CERN particle accelerator was (=) indeed the famous Higgs particle.** |
| **(48) Physicists agree that the CERN experiments did find a new particle that had never been seen before, but according to (e&eb) an international research team, there is (=) no conclusive evidence that the particle was indeed the Higgs particle.The research team has scrutinized the existing scientific data from CERN about the newfound particle and published their analysis in the journal Physical Review D. A member of this team is Mads Toudal Frandsen, associate professor at the Center for Cosmology and Particle Physics Phenomenology, Department of Physics, Chemistry and Pharmacy at the University of Southern Denmark. "The CERN data is generally taken as (=) evidence that the particle is the Higgs particle.** |

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| **(49) It is true that the Higgs particle can explain (eb) the data but there can be other explanations, we would also get this data from other particles", Mads Toudal Frandsen explains.** |
| **(50) The researchers' analysis does not (e) debunk the possibility that CERN has discovered the Higgs particle. That is still possible - but it is equally possible that it is a different kind of particle. "The current data is not precise enough to determine exactly what the particle is. It could be a number of other known particles", says Mads Toudal Frandsen. But if it wasn't the Higgs particle, which was found in CERN's particle accelerator, then what was it? "We believe that it may be a so-called techni-higgs particle. This particle is in some ways similar to the Higgs particle - hence half of the name", says Mads Toudal Frandsen.** |
| **(51) Although the techni-higgs particle and Higgs particle can easily be confused in experiments, they are two very different particles belonging to (e) two very different theories of how the universe was created.** |
| **(52) The Higgs particle is the missing piece in (eb) the theory called the Standard Model.** |
| **(53) This theory describes (eb) three of the four forces of nature. But it does not (e) explain what dark matter is - the substance that makes up (eb) most of the universe. A techni-higgs particle, if it exists, is (=) a completely different thing:** |
| **(54) "A techni-higgs particle is not (e) an elementary particle.** |
| **(55) Instead, it consists of so-called techni-quarks, which we believe are elementary. Techni-quarks may bind together in various ways to form (eb) for instance techni-higgs particles, while other combinations may form (eb) dark matter. Authors therefore expect to find several different particles at the LHC, all built by (e) techni-quarks", says Mads Toudal Frandsen.** |
| **(56) If techni-quarks exist, (eb) there must be a force to bind them together so that (e) they can form particles.** |
| **(57) None of the four known forces of nature (gravity, the electromagnetic force, the weak nuclear force and the strong nuclear force) are (=) any good at binding techni-quarks together.** |
| **(58) There must therefore be (=) a yet undiscovered force of nature. This force is called the the Technicolor force.** |
| **(59) What was found last year in CERN's accelerator could thus be (=) either the Higgs particle of the Standard Model or a light techni-higgs particle, composed of (e) two techni-quarks.** |
| **(60) Mads Toudal Frandsen believes that more data from CERN will probably be able to determine (eb) if it was a Higgs or a techni-higgs particle. If CERN gets an even more powerful accelerator, it will in principle be able to observe techni-quarks directly. More information: Technicolor Higgs boson in the light of LHC data. Phys. Rev. D 90, 035012th Alexander Belyaev, Matthew S. Brown, Roshan Foadi, and Mads T. Frandsen. journals.aps.org/prd/abstract/10.1103/PhysRevD.90.035012. On Arxiv: arxiv.org/abs/1309.2097 Journal reference: Physical Review D search and more info website: Provided by: University of Southern Denmark search and more info website** |

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| **The inflationary brane-antibrane universeClifford P. Burgess1,2, Mahbub Majumdar3, Detlef Nolte1, Fernando Quevedo3, Govindan Rajesh1 and Ren-Jie Zhang1 Published 13 August 2001 • Journal of High Energy Physics, Volume 2001, JHEP07(2001)** |
| **(61) Authors show how the motion through (e&eb) the extra dimensions of a gas of branes and antibranes can, under certain circumstances, produce (eb) an era of inflation as seen by (e) observers trapped on (e&eb) a 3-brane, with (e&eb) the inflaton being the inter-brane separation.** |
| **(62) Although most of discussion refers to (e) arbitrary ps, when we need to be specific they assume that they are (=) D-branes of type-II or type-I string theory.** |
| **(63) For realistic brane couplings, such as those arising in (eb) string theory, the inter-brane potentials are (=) too steep to inflate (eb) the universe for acceptably long times.** |
| **(64) However, for special regions of the parameter space of (e) brane-antibrane positions the brane motion is (=) slow enough for there to be (=) sufficient inflation.** |
| **(65) Inflation would be more generic in models where (e) the inter-brane interactions are much weaker.** |
| **(66) The spectrum of primordial density fluctuations predicted has (e) index n slightly less than 1, and acceptable amplitude, provided (e) that the extra dimensions have (e) linear size 1/r ~ 1012 GeV.** |
| **(67) Reheating occurs as in hybrid inflation, with (e&eb) the tachyonic instability of (e) the brane-antibrane system taking over for (e) small separations.** |
| **(68) The tachyon field can induce (eb) a cascade mechanism within which higher-dimension branes annihilate (e) into lower-dimension ones.** |
| **(69) They argue that such a cascade naturally stops with (e&eb) the production of 3-branes in (eb) 10-dimensional string theory.** |
| **Effect of the neutron lifetime on processes in the early universe A. V. Chechkin, A. V. Ivanchik, A. P. Serebrov, S. V. Bobashev First Online: 27 July 2016 DOI: 10.1134/S1063784216070069 Chechkin, A.V., Ivanchik, A.V., Serebrov, A.P. et al. Tech. Phys. (2016) 61: 1101. Doi: 10.1134/S1063784216070069** |
| **(70) The influence of the neutron lifetime on (e&eb) the abundance of light elements produced during (e&eb) the primordial nucleosynthesis early in (eb) the birth of the Universe is considered.** |
| **(71) Among light elements, namely, D, 3He, 4He, and 7Li, 4He proves to be (=) most sensitive to neutron lifetime τn.** |
| **(72) Astronomic data on the light element abundance also provide (eb) the best accuracy for 4He.** |
| **(73) The solution of a number of problems discussed in this paper requires (e) improving the accuracy of observations for the 4He abundance and refining the value of τn.** |

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| **Application of a Theory and Simulation based Convective Boundary Mixing model for AGB Star Evolution and Nucleosynthesis U. Battino, M. Pignatari, C. Ritter, F. Herwig, P. Denisenkov, J.W. Den Hartogh, R. Trappitsch, R. Hirschi, B. Freytag, F. Thielemann, B. Paxton (Submitted on 19 May 2016)** |
| **(74) The s-process nucleosynthesis in Asymptotic Giant Branch (AGB) stars depends on (e&eb) the modeling of (e) convective boundaries.** |
| **(75) Authors' present models and s-process simulations that adopt (e&eb) a treatment of convective boundaries based on (e) the results of hydrodynamic simulations and on (e&eb) the theory of mixing due to (e) gravity waves in (eb) the vicinity of convective boundaries.** |
| **(76) Hydrodynamics simulations suggest (eb) the presence of convective boundary mixing (CBM) at (eb) the bottom of the thermal pulse-driven convective zone.** |
| **(77) Similarly, convection-induced mixing processes are proposed for (e) the mixing below the convective envelope during (e&eb) third dredge-up where (e) the 13C pocket for (e) the s process in AGB stars forms.** |
| **(78) In this work apply apply a CBM model motivated by (e) simulations and theory to (e&eb) models with initial mass M = 2 and M = 3M?, and with (e&eb) initial metal content Z = 0:01 and Z = 0:02.** |
| **(79) As reported previously, the He-intershell abundance of 12C and 16O are increased by (e) CBM at the bottom of (e) pulse-driven convection zone.** |
| **(80) This mixing is affecting (e&eb) the 22Ne (alpha, n) 25Mg activation and (e&eb) the s-process effciency in the 13C-pocket.** |
| **(81) In our model CBM at the bottom of the convective envelope during (e&eb) the third dredgeup represents (eb) gravity wave mixing.** |
| **(82) Authors take further into account that hydrodynamic simulations indicate (eb) a declining mixing efficiency already about (e&eb) a pressure scale height from (e) the convective boundaries, compared to (e&eb) mixing-length theory.** |
| **(83) The final s-process abundances are characterized by (e) 0:36 < [s=Fe] < 0:78 and the heavy-to-light s-process ratio is (=) 0:23 < [hs=ls] < 0:45.** |
| **(84) Finally, authors compare results with (e&eb) stellar observations, pre-solar grain measurements and previous work. Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI:10.3847/0004-637X/827/1/30 Cite as: arXiv:1605.06159 [astro-ph.SR] (or arXiv:1605.06159v1 [astro-ph.SR] for this version)** |
| **MESA and NuGrid simulations of classical novae: CO and ONe nova nucleosynthesis Pavel A. Denissenkov (UVic), James W. Truran (Chicago), Marco Pignatari (Basel), Reto Trappitsch (Chicago), Christian Ritter (Frankfurt), Falk Herwig (UVic), Umberto Battino (Basel), Kiana Setoodehnia (McMaster), B. Paxton (UC Santa Barbara)** |

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| **(85) Classical novae are the result of (e) thermonuclear flashes of hydrogen accreted by (e) CO or ONe white dwarfs, leading (eb) eventually to the dynamic ejection of (e) the surface layers.** |
| **(86) These are observationally known to be (=) enriched in heavy elements, such as (=) C, O and Ne that must originate in (eb) layers below the H-flash convection zone.** |
| **(87) Building on their previous work, they now present stellar evolution simulations of (e) ONe novae and provide (eb) a comprehensive comparison of models with (e&eb) published ones.** |
| **(88) Some of authors models include exponential convective boundary mixing to account for (e) the observed enrichment of (e) the nova ejecta even when accreted material has (e) a solar abundance distribution.** |
| **(89) Models produce (eb) maximum temperature evolution profiles and nucleosynthesis yields (eb) in good agreement with (e&eb) models that generate (eb) enriched ejecta by assuming (eb) that the accreted material was pre-mixed.** |
| **(90) Authors confirm for (e) ONe novae the result we reported previously, i.e.\ we found that 3He could be produced {\it in situ} in (eb) solar-composition envelopes accreted with (e&eb) slow rates ($M˙$<10−10$M_\odot$/yr) by (e) cold (TWD<107 K) CO WDs, and that convection was triggered by (eb) 3He burning before (e) the nova outburst in that case.** |
| **(91) In addition, authors now find that the interplay between the 3He production and (e&eb) destruction in the solar-composition envelope accreted with (e&eb) an intermediate rate, e.g.\ $M˙$=10−10$M_\odot$/yr, by the 1.15$M_\odot$ ONe WD with (e&eb) a relatively high initial central temperature, e.g.\ TWD=15×106 K, leads to (eb) the formation of a thick radiative buffer zone that separates (e&eb) the bottom of the convective envelope from (e) the WD surface. (Abridged) Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI: 10.1093/mnras/stu1000 Cite as: arXiv:1303.6265 [astro-ph.SR] (or arXiv:1303.6265v2 [astro-ph.SR] for this version)** |
| **Hybrid C-O-Ne White Dwarfs as Progenitors of Diverse SNe Ia Pavel Denissenkov (UVic), James Truran (U. Chicago), Falk Herwig (UVic), Sam Jones (UVic), Bill Paxton (UC Santa Barbara), Ken'ichi Nomoto (U. Tokyo), Toshio Suzuki (Nihon U.), Hiroshi Toki (Osaka U.)(Submitted on 6 Nov 2014)** |
| **(92) When carbon is ignited off-center in a CO core of a super-AGB star (eb) its burning in a convective shell tends to propagate to (e&eb) the center.** |
| **(93) Whether the C flame will actually be able to reach (eb) the center depends on (e&eb) the efficiency of (e) extra mixing beneath (e&eb) the C convective shell.** |
| **(94) Whereas thermohaline mixing is too inefficient to interfere with the C-flame propagation, (eb) convective boundary mixing can prevent (e) the C burning from reaching (eb) the center.** |
| **(95) As a result, a C-O-Ne white dwarf (WD) is (=) formed, after the star has lost (e) its envelope.** |
| **(96) Such a "hybrid" WD has (e) a small CO core surrounded by (e) a thick ONe zone.** |

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| **(97) In the proposed 1D stellar evolution computations the hybrid WD is allowed to (eb) accrete C-rich material, as if it were in a close binary system and (e&eb) accreted H-rich material from (e) its companion with (e&eb) a sufficiently high rate at which the accreted H would be processed into (e&eb) He under (e&eb) stationary conditions, assuming (eb) that He could then be transformed into (e&eb) C.** |
| **(98) When the mass of the accreting WD approaches (eb) the Chandrasekhar limit, they find a series of convective Urca shell flashes associated with (e&eb) high abundances of 23Na and 25Mg.** |
| **(99) They are followed by off-center C ignition leading to (eb) convection that occupies (e) almost the entire star.** |
| **(100) To model the Urca processes, authors use the most recent well-resolved data for (e) their reaction and (e&eb) neutrino-energy loss rates.** |
| **(101) Because of the emphasized uncertainty of the convective Urca process in hybrid WD models of (e) SN Ia progenitors, authors consider a number of their potentially possible alternative instances for (e) different mixing assumptions, all of which reach (eb) a phase of explosive C ignition, either off or in (eb) the center.** |
| **(102) Our hybrid SN Ia progenitor models have (e) much lower C to O abundance ratios at (eb) the moment of the explosive C ignition than (e) their pure CO counterparts, which may explain (eb) the observed diversity of the SNe Ia. subjects: Solar and Stellar Astrophysics (astro-ph.SR)cite as: arXiv:1411.1471 [astro-ph.SR](or arXiv:1411.1471v1 [astro-ph.SR] for this version)** |
| **Rejecting proposed dense-matter equations of state with quiescent low-mass X-ray binaries Sebastien Guillot, Robert E. Rutledge (Submitted on 15 Sep 2014 (v1), last revised 20 Oct 2014 (this version, v2))** |
| **(103) Neutrons stars are unique laboratories to (e) discriminate between the various proposed equations of (e) state of matter at and above nuclear density.** |
| **(104) One sub-class of neutron stars - those inside (eb) quiescent low-mass X-ray binaries (qLMXBs) - produce (eb) a thermal surface emission from (e) which the neutron star radius (R_NS) can be measured, using (e) the widely accepted observational scenario for qLMXBs, assuming (eb) unmagnetized H atmospheres.** |
| **(105) In a combined spectral analysis, this work first reproduces (eb) a previously published measurement of the \rns, assumed to be (=) the same for all neutron stars, using (e) a slightly expanded data set.** |
| **(106) The radius measured is R_NS = 9.4 +/-1.2 km. On the basis of spectral analysis alone, this measured value is not affected by imposing (e&eb) an assumption of causality in the core.** |
| **(107) However, the assumptions underlying this R_NS measurement would be falsified by (e) the observation of any neutron star with (e&eb) a mass >2.6 Msun, since radii <11 km would be rejected if (e) causality is assumed, which would exclude (eb) most of the R_NS parameter space obtained in this analysis.** |

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| **(108) Finally, this work directly tests (e&eb) a selection of dense matter equations of states: WFF1, AP4, MPA1, PAL1, MS0, and three versions of equations of state produced through (e&eb) chiral effective theory.** |
| **(109) Two of those, MSO and PAL1, are rejected at (eb) the 99% confidence level, accounting for (e) all quantifiable uncertainties, while (e) the other cannot be excluded at >99% certainty. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE) DOI: 10.1088/2041-8205/796/1/L3 Cite as: arXiv:1409.4306 [astro-ph.HE] (or arXiv:1409.4306v2 [astro-ph.HE] for this version** |

## SECTION TWO

## Scattered Photon Cannot Perform A "Strong" Measurement Of The Particles Position

**INTRODUCTION—VARIABLES USED**

**Viewpoint: Measuring Quantum Kicks from a Beam of Light Amir Safavi-Naeini, 348 Via Pueblo Mall, Stanford University, Stanford, CA 94305-4088, USA June 13, 2016• Physics 9, 64 Force sensors levitated by light have reached the quantum regime, in which their sensitivity is limited by the momentum kicks of individual photons. Figure caption expand figure APS/Alan Stonebraker Figures and some sentences are deleted due to spatial restrictions Please excuse me**

(1) For the experimental parameters chosen by Novotny's group, the nanosphere's zero-point fluctuation length is (=) approximately 10 pm.
(2) This is the length scale on which a photon must measure (eb) the particle's position in order to discern (e&eb) quantum fluctuations.
(3) However, the photons from the trapping laser have a wavelength of 1 μ$\mu$m, which is one hundred thousand times larger than (e) the fluctuation length.
(4) As such, a scattered photon cannot (e) perform a "strong" measurement of (e) the particles position.
(5) In other words, the quantum back action from (e) the scattering of a single photon can have (e) only a minor effect on (e&eb) the dynamics of the particle.
(6) However, the cumulative momentum kicks from (e) many scattering events can change (e&eb) the state of the nanosphere.

**NOTATION**

**Module One**

. For the experimental parameters chosen by Novotny's group, the nanosphere's zero-point fluctuation length is (=) approximately 10 pm

$G_{13}$ : Category one of experimental parameters chosen by Novotny's group, the nanosphere's zero-point fluctuation length

$G_{14}$ : Category two of SAS

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| $G_{15}$ : Category three of SAS | |
| $T_{13}$ : Category one of approximately 10 pm<br>$T_{14}$ : Category two of SAS<br>$T_{15}$ : Category three of SAS | |
| **Module Two**<br>This is the length scale on which a photon must measure (eb) the particle's position in order to discern (e&eb) quantum fluctuations | |
| $G_{16}$ : Category one of length scale on which a photon must measure<br>$G_{17}$ : Category two of SAS<br>$G_{18}$ : Category three of SAS | |
| $T_{16}$ : Category one of particle's position in order to discern (e&eb) quantum fluctuations<br>$T_{17}$ : Category two of SAS<br>$T_{18}$ : Category three of SAS | |
| **Module three**<br>This is the **length scale on which a photon must measure the particle's position in order to discern** (e&eb) quantum fluctuations | |
| $G_{20}$ : Category one of **length scale on which a photon must measure the particle's position in order to discern;** quantum fluctuations<br>$G_{21}$ : Category two of SAS<br>$G_{22}$ : Category three of SAS | |
| $T_{20}$ : Category one of quantum fluctuations **;length scale on which a photon must measure the particle's position in order to discern**<br>$T_{21}$ : Category two of SAS<br>$T_{22}$ : Category three of SAS | |
| **Module four**<br>However, the photons from the trapping laser have a wavelength of 1 μ$\mu$m, which is one hundred thousand times larger than (e) the fluctuation length | |
| $G_{24}$ : Category one of fluctuation length | |

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| $G_{25}$ : Category two of SAS<br><br>$G_{26}$ : Category three of SAS | |
| $T_{24}$ : Category one of photons from the trapping laser have a wavelength of 1 μ$\mu$m, which is one hundred thousand times larger<br><br>$T_{25}$ : Category two of SAS<br><br>$T_{26}$ : Category three of SAS | |
| **Module five**<br><br>As such, a scattered photon cannot perform a "strong" measurement of (e) the particles position | |
| $G_{28}$ : Category one of particles position<br><br>$G_{29}$ : Category two of SAS<br><br>$G_{30}$ : Category three of SAS | |
| $T_{28}$ : Category one of scattered photon cannot perform a "strong" measurement<br><br>$T_{29}$ : Category two of SAS<br><br>$T_{30}$ : Category three of SAS | |
| **Module six**<br><br>In other words, the quantum back action from (e) the scattering of a single photon can have (e) only a minor effect on (e&eb) the dynamics of the particle | |
| $G_{32}$ : Category one of scattering of a single photon can have (e) only a minor effect on (e&eb) the dynamics of the particle<br><br>$G_{33}$ : Category two of SAS<br><br>$G_{34}$ : Category three of SAS | |
| $T_{32}$ : Category one of quantum back action<br><br>$T_{33}$ : Category two of SAS<br><br>$T_{34}$ : Category three of SAS | |
| **Module seven**<br><br>other words, the quantum back action from the scattering of a single photon can have (e) only a minor effect on (e&eb) the dynamics of the particle | |
| $G_{36}$ : Category one of only a minor effect on (e&eb) the dynamics of the particle | |

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| $G_{37}$ : Category two of SAS<br><br>$G_{38}$ : Category three of SAS | |
| $T_{36}$ : Category one of quantum back action from the scattering of a single photon<br><br>$T_{37}$ : Category two of SAS<br><br>$T_{38}$ : Category three of SAS | |
| **Module eight**<br><br>other words, the quantum back action from the scattering of a single photon can have only a minor effect on (e&eb) the dynamics of the particle | |
| $G_{40}$ : Category one of **quantum back action from the scattering of a single photon can have only a minor effect;** dynamics of the particle<br><br>$G_{41}$ : Category two of SAS<br><br>$G_{42}$ : Category three of SAS | |
| $T_{40}$ : Category one of dynamics of the particle; **quantum back action from the scattering of a single photon can have only a minor effect**<br><br>$T_{41}$ : Category two of SAS<br><br>$T_{42}$ : Category three of SAS | |
| **Module Nine**<br><br>However, the cumulative momentum kicks from (e) many scattering events can change (e&eb) the state of the nanosphere | |
| $G_{44}$ : Category one of many scattering events can change (e&eb) the state of the nanosphere<br><br>$G_{45}$ : Category two of SAS<br><br>$G_{46}$ : Category three of SAS | |
| $T_{44}$ : Category one of cumulative momentum kicks<br><br>$T_{45}$ : Category two of SAS<br><br>$T_{46}$ : Category three of SAS | |
| **The Coefficients:** | |
| $(a_{13})^{(1)}, (a_{14})^{(1)}, (a_{15})^{(1)}, (b_{13})^{(1)}, (b_{14})^{(1)}, (b_{15})^{(1)}$ $(a_{16})^{(2)}, (a_{17})^{(2)}, (a_{18})^{(2)}$ $(b_{16})^{(2)}, (b_{17})^{(2)}, (b_{18})^{(2)}$:<br>$(a_{20})^{(3)}, (a_{21})^{(3)}, (a_{22})^{(3)}$ , $(b_{20})^{(3)}, (b_{21})^{(3)}, (b_{22})^{(3)}$<br>$(a_{24})^{(4)}, (a_{25})^{(4)}, (a_{26})^{(4)}, (b_{24})^{(4)}, (b_{25})^{(4)}, (b_{26})^{(4)}, (b_{28})^{(5)}, (b_{29})^{(5)}, (b_{30})^{(5)}, (a_{28})^{(5)}, (a_{29})^{(5)}, (a_{30})^{(5)},$ | |

| | |
|---|---|
| $(a_{32})^{(6)}, (a_{33})^{(6)}, (a_{34})^{(6)}, (b_{32})^{(6)}, (b_{33})^{(6)}, (b_{34})^{(6)}$<br>$(a_{36})^{(7)}, (a_{37})^{(7)}, (a_{38})^{(7)}, (b_{36})^{(7)}, (b_{37})^{(7)}, (b_{38})^{(7)}$<br>$(a_{40})^{(8)}, (a_{41})^{(8)}, (a_{42})^{(8)}, (b_{40})^{(8)}, (b_{41})^{(8)}, (b_{42})^{(8)}$<br>$(a_{44})^{(9)}, (a_{45})^{(9)}, (a_{46})^{(9)}, (b_{44})^{(9)}, (b_{45})^{(9)}, (b_{46})^{(9)}$<br><br>are Accentuation coefficients<br>$(a'_{13})^{(1)}, (a'_{14})^{(1)}, (a'_{15})^{(1)}, (b'_{13})^{(1)}, (b'_{14})^{(1)}, (b'_{15})^{(1)}, (a'_{16})^{(2)}, (a'_{17})^{(2)}, (a'_{18})^{(2)}, \ (b'_{16})^{(2)}, (b'_{17})^{(2)}, (b'_{18})^{(2)}$<br>$, (a'_{20})^{(3)}, (a'_{21})^{(3)}, (a'_{22})^{(3)}, (b'_{20})^{(3)}, (b'_{21})^{(3)}, (b'_{22})^{(3)}$<br>$(a'_{24})^{(4)}, (a'_{25})^{(4)}, (a'_{26})^{(4)}, (b'_{24})^{(4)}, (b'_{25})^{(4)}, (b'_{26})^{(4)}, (b'_{28})^{(5)}, (b'_{29})^{(5)}, (b'_{30})^{(5)}\ (a'_{28})^{(5)}, (a'_{29})^{(5)}, (a'_{30})^{(5)}$ ,<br>$(a'_{32})^{(6)}, (a'_{33})^{(6)}, (a'_{34})^{(6)}, (b'_{32})^{(6)}, (b'_{33})^{(6)}, (b'_{34})^{(6)}$<br>$(a'_{36})^{(7)}, (a'_{37})^{(7)}, (a'_{38})^{(7)}, (b'_{36})^{(7)}, (b'_{37})^{(7)}, (b'_{38})^{(7)}$,<br>$(a'_{40})^{(8)}, (a'_{41})^{(8)}, (a'_{42})^{(8)}, (b'_{40})^{(8)}, (b'_{41})^{(8)}, (b'_{42})^{(8)}$,<br>$(a'_{44})^{(9)}, (a'_{45})^{(9)}, (a'_{46})^{(9)}, (b'_{44})^{(9)}, (b'_{45})^{(9)}, (b'_{46})^{(9)}$,<br>are Dissipation coefficients | |
| **Module Numbered One**<br><br>**The differential system of this model is now (Module Numbered one)** | |
| $\frac{dG_{13}}{dt} = (a_{13})^{(1)}G_{14} - \left[(a'_{13})^{(1)} + (a''_{13})^{(1)}(T_{14}, t)\right]G_{13}$ | 1 |
| $\frac{dG_{14}}{dt} = (a_{14})^{(1)}G_{13} - \left[(a'_{14})^{(1)} + (a''_{14})^{(1)}(T_{14}, t)\right]G_{14}$ | 2 |
| $\frac{dG_{15}}{dt} = (a_{15})^{(1)}G_{14} - \left[(a'_{15})^{(1)} + (a''_{15})^{(1)}(T_{14}, t)\right]G_{15}$ | 3 |
| $\frac{dT_{13}}{dt} = (b_{13})^{(1)}T_{14} - \left[(b'_{13})^{(1)} - (b''_{13})^{(1)}(G, t)\right]T_{13}$ | 4 |
| $\frac{dT_{14}}{dt} = (b_{14})^{(1)}T_{13} - \left[(b'_{14})^{(1)} - (b''_{14})^{(1)}(G, t)\right]T_{14}$ | 5 |
| $\frac{dT_{15}}{dt} = (b_{15})^{(1)}T_{14} - \left[(b'_{15})^{(1)} - (b''_{15})^{(1)}(G, t)\right]T_{15}$ | 6 |
| $+(a''_{13})^{(1)}(T_{14}, t) =$ First augmentation factor | |
| $-(b''_{13})^{(1)}(G, t) =$ First detritions factor | |
| **Module Numbered Two**<br><br>**The differential system of this model is now ( Module numbered two)** | |
| $\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \left[(a'_{16})^{(2)} + (a''_{16})^{(2)}(T_{17}, t)\right]G_{16}$ | 7 |
| $\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \left[(a'_{17})^{(2)} + (a''_{17})^{(2)}(T_{17}, t)\right]G_{17}$ | 8 |
| $\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \left[(a'_{18})^{(2)} + (a''_{18})^{(2)}(T_{17}, t)\right]G_{18}$ | 9 |
| $\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \left[(b'_{16})^{(2)} - (b''_{16})^{(2)}\left((G_{19}), t\right)\right]T_{16}$ | 10 |
| $\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \left[(b'_{17})^{(2)} - (b''_{17})^{(2)}\left((G_{19}), t\right)\right]T_{17}$ | 11 |
| $\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \left[(b'_{18})^{(2)} - (b''_{18})^{(2)}\left((G_{19}), t\right)\right]T_{18}$ | 12 |
| $+(a''_{16})^{(2)}(T_{17}, t) =$ First augmentation factor | |
| $-(b''_{16})^{(2)}\left((G_{19}), t\right) =$ First detritions factor | |
| **Module Numbered Three**<br><br>**The differential system of this model is now (Module numbered three)** | |

| | |
|---|---|
| $\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \left[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21},t)\right]G_{20}$ | 13 |
| $\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \left[(a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)}(T_{21},t)\right]G_{21}$ | 14 |
| $\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \left[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21},t)\right]G_{22}$ | 15 |
| $\frac{dT_{20}}{dt} = (b_{20})^{(3)}T_{21} - \left[(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23},t)\right]T_{20}$ | 16 |
| $\frac{dT_{21}}{dt} = (b_{21})^{(3)}T_{20} - \left[(b_{21}^{'})^{(3)} - (b_{21}^{''})^{(3)}(G_{23},t)\right]T_{21}$ | 17 |
| $\frac{dT_{22}}{dt} = (b_{22})^{(3)}T_{21} - \left[(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23},t)\right]T_{22}$ | 18 |
| $+(a_{20}^{''})^{(3)}(T_{21},t) =$ First augmentation factor | |
| $-(b_{20}^{''})^{(3)}(G_{23},t) =$ First detritions factor | |
| **Module Numbered Four**<br><br>**The differential system of this model is now (Module numbered Four)** | |
| $\frac{dG_{24}}{dt} = (a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25},t)\right]G_{24}$ | 19 |
| $\frac{dG_{25}}{dt} = (a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25},t)\right]G_{25}$ | 20 |
| $\frac{dG_{26}}{dt} = (a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25},t)\right]G_{26}$ | 21 |
| $\frac{dT_{24}}{dt} = (b_{24})^{(4)}T_{25} - \left[(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}((G_{27}),t)\right]T_{24}$ | 22 |
| $\frac{dT_{25}}{dt} = (b_{25})^{(4)}T_{24} - \left[(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}((G_{27}),t)\right]T_{25}$ | 23 |
| $\frac{dT_{26}}{dt} = (b_{26})^{(4)}T_{25} - \left[(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}((G_{27}),t)\right]T_{26}$ | 24 |
| $+(a_{24}^{''})^{(4)}(T_{25},t) =$ **First augmentation factor** | |
| $-(b_{24}^{''})^{(4)}((G_{27}),t) =$ **First detritions factor** | |
| **Module Numbered Five**:<br><br>**The differential system of this model is now (Module number five)** | |
| $\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29},t)\right]G_{28}$ | 25 |
| $\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29},t)\right]G_{29}$ | 26 |
| $\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \left[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29},t)\right]G_{30}$ | 27 |
| $\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \left[(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}((G_{31}),t)\right]T_{28}$ | 28 |
| $\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \left[(b_{29}^{'})^{(5)} - (b_{29}^{''})^{(5)}((G_{31}),t)\right]T_{29}$ | 29 |
| $\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \left[(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}((G_{31}),t)\right]T_{30}$ | 30 |
| $+(a_{28}^{''})^{(5)}(T_{29},t) =$ **First augmentation factor** | |
| $-(b_{28}^{''})^{(5)}((G_{31}),t) =$ **First detritions factor** | |
| **Module Numbered Six**<br><br>**The differential system of this model is now (Module numbered Six)** | |
| $\frac{dG_{32}}{dt} = (a_{32})^{(6)}G_{33} - \left[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33},t)\right]G_{32}$ | 31 |
| $\frac{dG_{33}}{dt} = (a_{33})^{(6)}G_{32} - \left[(a_{33}^{'})^{(6)} + (a_{33}^{''})^{(6)}(T_{33},t)\right]G_{33}$ | 32 |
| $\frac{dG_{34}}{dt} = (a_{34})^{(6)}G_{33} - \left[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33},t)\right]G_{34}$ | 33 |

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|---|---|
| $\frac{dT_{32}}{dt} = (b_{32})^{(6)}T_{33} - \left[(b'_{32})^{(6)} - (b''_{32})^{(6)}((G_{35}),t)\right]T_{32}$ | 34 |
| $\frac{dT_{33}}{dt} = (b_{33})^{(6)}T_{32} - \left[(b'_{33})^{(6)} - (b''_{33})^{(6)}((G_{35}),t)\right]T_{33}$ | 35 |
| $\frac{dT_{34}}{dt} = (b_{34})^{(6)}T_{33} - \left[(b'_{34})^{(6)} - (b''_{34})^{(6)}((G_{35}),t)\right]T_{34}$ | 36 |
| $+(a''_{32})^{(6)}(T_{33},t) =$ **First augmentation factor** | |
| **Module Numbered Seven:**<br><br>**The differential system of this model is now (Seventh Module)** | |
| $\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \left[(a'_{36})^{(7)} + (a''_{36})^{(7)}(T_{37},t)\right]G_{36}$ | 37 |
| $\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \left[(a'_{37})^{(7)} + (a''_{37})^{(7)}(T_{37},t)\right]G_{37}$ | *38* |
| $\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \left[(a'_{38})^{(7)} + (a''_{38})^{(7)}(T_{37},t)\right]G_{38}$ | *39* |
| $\frac{dT_{36}}{dt} = (b_{36})^{(7)}T_{37} - \left[(b'_{36})^{(7)} - (b''_{36})^{(7)}((G_{39}),t)\right]T_{36}$ | *40* |
| $\frac{dT_{37}}{dt} = (b_{37})^{(7)}T_{36} - \left[(b'_{37})^{(7)} - (b''_{37})^{(7)}((G_{39}),t)\right]T_{37}$ | *41* |
| $\frac{dT_{38}}{dt} = (b_{38})^{(7)}T_{37} - \left[(b'_{38})^{(7)} - (b''_{38})^{(7)}((G_{39}),t)\right]T_{38}$ | *42* |
| $+(a''_{36})^{(7)}(T_{37},t) =$ **First augmentation factor** | |
| **Module Numbered Eight**<br><br>The differential system of this model is now | |
| $\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \left[(a'_{40})^{(8)} + (a''_{40})^{(8)}(T_{41},t)\right]G_{40}$ | *43* |
| $\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \left[(a'_{41})^{(8)} + (a''_{41})^{(8)}(T_{41},t)\right]G_{41}$ | *44* |
| $\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \left[(a'_{42})^{(8)} + (a''_{42})^{(8)}(T_{41},t)\right]G_{42}$ | *45* |
| $\frac{dT_{40}}{dt} = (b_{40})^{(8)}T_{41} - \left[(b'_{40})^{(8)} - (b''_{40})^{(8)}((G_{43}),t)\right]T_{40}$ | *46* |
| $\frac{dT_{41}}{dt} = (b_{41})^{(8)}T_{40} - \left[(b'_{41})^{(8)} - (b''_{41})^{(8)}((G_{43}),t)\right]T_{41}$ | *47* |
| $\frac{dT_{42}}{dt} = (b_{42})^{(8)}T_{41} - \left[(b'_{42})^{(8)} - (b''_{42})^{(8)}((G_{43}),t)\right]T_{42}$ | *48* |
| **Module Numbered Nine**<br><br>The differential system of this model is now | |
| $\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \left[(a'_{44})^{(9)} + (a''_{44})^{(9)}(T_{45},t)\right]G_{44}$ | 49 |
| $\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \left[(a'_{45})^{(9)} + (a''_{45})^{(9)}(T_{45},t)\right]G_{45}$ | 50 |
| $\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \left[(a'_{46})^{(9)} + (a''_{46})^{(9)}(T_{45},t)\right]G_{46}$ | 51 |
| $\frac{dT_{44}}{dt} = (b_{44})^{(9)}T_{45} - \left[(b'_{44})^{(9)} - (b''_{44})^{(9)}((G_{47}),t)\right]T_{44}$ | 52 |
| $\frac{dT_{45}}{dt} = (b_{45})^{(9)}T_{44} - \left[(b'_{45})^{(9)} - (b''_{45})^{(9)}((G_{47}),t)\right]T_{45}$ | 53 |
| $\frac{dT_{46}}{dt} = (b_{46})^{(9)}T_{45} - \left[(b'_{46})^{(9)} - (b''_{46})^{(9)}((G_{47}),t)\right]T_{46}$ | 54 |
| $+(a''_{44})^{(9)}(T_{45},t) =$ **First augmentation factor** | |
| $-(b''_{44})^{(9)}((G_{47}),t) =$ **First detrition factor** | |

$$\frac{dG_{13}}{dt} = (a_{13})^{(1)}G_{14} - \left[\begin{array}{ccc} (a_{13}^{'})^{(1)}\boxed{+(a_{13}^{''})^{(1)}(T_{14},t)} & \boxed{+(a_{16}^{''})^{(2,2,)}(T_{17},t)} & \boxed{+(a_{20}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{24}^{''})^{(4,4,4,4,)}(T_{25},t)} & \boxed{+(a_{28}^{''})^{(5,5,5,5,)}(T_{29},t)} & \boxed{+(a_{32}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{36}^{''})^{(7,7)}(T_{37},t)} & \boxed{+(a_{40}^{''})^{(8,8)}(T_{41},t)} & \boxed{+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{array}\right] G_{13} \quad (55)$$

$$\frac{dG_{14}}{dt} = (a_{14})^{(1)}G_{13} - \left[\begin{array}{ccc} (a_{14}^{'})^{(1)}\boxed{+(a_{14}^{''})^{(1)}(T_{14},t)} & \boxed{+(a_{17}^{''})^{(2,2,)}(T_{17},t)} & \boxed{+(a_{21}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{25}^{''})^{(4,4,4,4,)}(T_{25},t)} & \boxed{+(a_{29}^{''})^{(5,5,5,5,)}(T_{29},t)} & \boxed{+(a_{33}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{37}^{''})^{(7,7)}(T_{37},t)} & \boxed{+(a_{41}^{''})^{(8,8)}(T_{41},t)} & \boxed{+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{array}\right] G_{14} \quad (56)$$

$$\frac{dG_{15}}{dt} = (a_{15})^{(1)}G_{14} - \left[\begin{array}{ccc} (a_{15}^{'})^{(1)}\boxed{+(a_{15}^{''})^{(1)}(T_{14},t)} & \boxed{+(a_{18}^{''})^{(2,2,)}(T_{17},t)} & \boxed{+(a_{22}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{26}^{''})^{(4,4,4,4,)}(T_{25},t)} & \boxed{+(a_{30}^{''})^{(5,5,5,5,)}(T_{29},t)} & \boxed{+(a_{34}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{38}^{''})^{(7,7)}(T_{37},t)} & \boxed{+(a_{42}^{''})^{(8,8)}(T_{41},t)} & \boxed{+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{array}\right] G_{15} \quad (57)$$

Where $\boxed{(a_{13}^{''})^{(1)}(T_{14},t)}$, $\boxed{(a_{14}^{''})^{(1)}(T_{14},t)}$, $\boxed{(a_{15}^{''})^{(1)}(T_{14},t)}$ are first augmentation coefficients for category 1, 2 and 3

$\boxed{+(a_{16}^{''})^{(2,2,)}(T_{17},t)}$, $\boxed{+(a_{17}^{''})^{(2,2,)}(T_{17},t)}$, $\boxed{+(a_{18}^{''})^{(2,2,)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{20}^{''})^{(3,3,)}(T_{21},t)}$, $\boxed{+(a_{21}^{''})^{(3,3,)}(T_{21},t)}$, $\boxed{+(a_{22}^{''})^{(3,3,)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{24}^{''})^{(4,4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{25}^{''})^{(4,4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{26}^{''})^{(4,4,4,4,)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{28}^{''})^{(5,5,5,5,)}(T_{29},t)}$, $\boxed{+(a_{29}^{''})^{(5,5,5,5,)}(T_{29},t)}$, $\boxed{+(a_{30}^{''})^{(5,5,5,5,)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{32}^{''})^{(6,6,6,6,)}(T_{33},t)}$, $\boxed{+(a_{33}^{''})^{(6,6,6,6,)}(T_{33},t)}$, $\boxed{+(a_{34}^{''})^{(6,6,6,6,)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{38}^{''})^{(7,7)}(T_{37},t)}$ $\boxed{+(a_{37}^{''})^{(7,7)}(T_{37},t)}$ $\boxed{+(a_{36}^{''})^{(7,7)}(T_{37},t)}$ are seventh augmentation coefficient for 1,2,3

$\boxed{+(a_{40}^{''})^{(8,8)}(T_{41},t)}$ $\boxed{+(a_{41}^{''})^{(8,8)}(T_{41},t)}$ $\boxed{+(a_{42}^{''})^{(8,8)}(T_{41},t)}$ are eight augmentation coefficient for 1,2,3

$\boxed{+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)}$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{13}}{dt} = (b_{13})^{(1)}T_{14} - \left[\begin{array}{ccc} (b_{13}^{'})^{(1)}\boxed{-(b_{13}^{''})^{(1)}(G,t)} & \boxed{-(b_{16}^{''})^{(2,2,)}(G_{19},t)} & \boxed{-(b_{20}^{''})^{(3,3,)}(G_{23},t)} \\ \boxed{-(b_{24}^{''})^{(4,4,4,4,)}(G_{27},t)} & \boxed{-(b_{28}^{''})^{(5,5,5,5,)}(G_{31},t)} & \boxed{-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)} \\ \boxed{-(b_{36}^{''})^{(7,7,)}(G_{39},t)} & \boxed{-(b_{40}^{''})^{(8,8)}(G_{43},t)} & \boxed{-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)} \end{array}\right] T_{13} \quad (58)$$

$$\frac{dT_{14}}{dt} = (b_{14})^{(1)}T_{13} - \left[\begin{array}{ccc} (b_{14}^{'})^{(1)}\boxed{-(b_{14}^{''})^{(1)}(G,t)} & \boxed{-(b_{17}^{''})^{(2,2,)}(G_{19},t)} & \boxed{-(b_{21}^{''})^{(3,3,)}(G_{23},t)} \\ \boxed{-(b_{25}^{''})^{(4,4,4,4,)}(G_{27},t)} & \boxed{-(b_{29}^{''})^{(5,5,5,5,)}(G_{31},t)} & \boxed{-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)} \\ \boxed{-(b_{37}^{''})^{(7,7,)}(G_{39},t)} & \boxed{-(b_{41}^{''})^{(8,8)}(G_{43},t)} & \boxed{-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)} \end{array}\right] T_{14} \quad (59)$$

$$\frac{dT_{15}}{dt} = (b_{15})^{(1)}T_{14} - \begin{bmatrix} (b_{15}^{'})^{(1)} & -(b_{15}^{''})^{(1)}(G,t) & -(b_{18}^{''})^{(2,2,)}(G_{19},t) & -(b_{22}^{''})^{(3,3,)}(G_{23},t) \\ -(b_{26}^{''})^{(4,4,4,4,)}(G_{27},t) & -(b_{30}^{''})^{(5,5,5,5,)}(G_{31},t) & -(b_{34}^{''})^{(6,6,6,6,)}(G_{35},t) \\ -(b_{38}^{''})^{(7,7,)}(G_{39},t) & -(b_{42}^{''})^{(8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{15} \quad (60)$$

Where $-(b_{13}^{''})^{(1)}(G,t)$, $-(b_{14}^{''})^{(1)}(G,t)$, $-(b_{15}^{''})^{(1)}(G,t)$ are first detrition coefficients for category 1, 2 and 3

$-(b_{16}^{''})^{(2,2,)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3

$-(b_{20}^{''})^{(3,3,)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3

$-(b_{24}^{''})^{(4,4,4,4,)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3

$-(b_{28}^{''})^{(5,5,5,5,)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3

$-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,)}(G_{35},t)$ are sixth detrition coefficients for category 1, 2 and 3

$-(b_{37}^{''})^{(7,7,)}(G_{39},t)$, $-(b_{36}^{''})^{(7,7,)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3

$-(b_{40}^{''})^{(8,8)}(G_{43},t)$ $-(b_{41}^{''})^{(8,8)}(G_{43},t)$ $-(b_{42}^{''})^{(8,8)}(G_{43},t)$ are eight detrition coefficients for category 1, 2 and 3

$-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3

$$\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \begin{bmatrix} (a_{16}^{'})^{(2)} & +(a_{16}^{''})^{(2)}(T_{17},t) & +(a_{13}^{''})^{(1,1,)}(T_{14},t) & +(a_{20}^{''})^{(3,3,3)}(T_{21},t) \\ +(a_{24}^{''})^{(4,4,4,4,4)}(T_{25},t) & +(a_{28}^{''})^{(5,5,5,5,5)}(T_{29},t) & +(a_{32}^{''})^{(6,6,6,6,6)}(T_{33},t) \\ +(a_{36}^{''})^{(7,7,7)}(T_{37},t) & +(a_{40}^{''})^{(8,8,8)}(T_{41},t) & +(a_{44}^{''})^{(9,9)}(T_{45},t) \end{bmatrix} G_{16} \quad (61)$$

$$\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \begin{bmatrix} (a_{17}^{'})^{(2)} & +(a_{17}^{''})^{(2)}(T_{17},t) & +(a_{14}^{''})^{(1,1,)}(T_{14},t) & +(a_{21}^{''})^{(3,3,3)}(T_{21},t) \\ +(a_{25}^{''})^{(4,4,4,4,4)}(T_{25},t) & +(a_{29}^{''})^{(5,5,5,5,5)}(T_{29},t) & +(a_{33}^{''})^{(6,6,6,6,6)}(T_{33},t) \\ +(a_{37}^{''})^{(7,7,7)}(T_{37},t) & +(a_{41}^{''})^{(8,8,8)}(T_{41},t) & +(a_{45}^{''})^{(9,9)}(T_{45},t) \end{bmatrix} G_{17} \quad (62)$$

$$\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \begin{bmatrix} (a_{18}^{'})^{(2)} & +(a_{18}^{''})^{(2)}(T_{17},t) & +(a_{15}^{''})^{(1,1,)}(T_{14},t) & +(a_{22}^{''})^{(3,3,3)}(T_{21},t) \\ +(a_{26}^{''})^{(4,4,4,4,4)}(T_{25},t) & +(a_{30}^{''})^{(5,5,5,5,5)}(T_{29},t) & +(a_{34}^{''})^{(6,6,6,6,6)}(T_{33},t) \\ +(a_{38}^{''})^{(7,7,7)}(T_{37},t) & +(a_{42}^{''})^{(8,8,8)}(T_{41},t) & +(a_{46}^{''})^{(9,9)}(T_{45},t) \end{bmatrix} G_{18} \quad (63)$$

Where $+(a_{16}^{''})^{(2)}(T_{17},t)$, $+(a_{17}^{''})^{(2)}(T_{17},t)$, $+(a_{18}^{''})^{(2)}(T_{17},t)$ are first augmentation coefficients for category 1, 2 and 3

$+(a_{13}^{''})^{(1,1,)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,)}(T_{14},t)$ are second augmentation coefficient for category 1, 2 and 3

$+(a_{20}^{''})^{(3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3)}(T_{21},t)$ are third augmentation coefficient for

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| category 1, 2 and 3 $+(a_{24}^{''})^{(4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3 $+(a_{28}^{''})^{(5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficient for category 1, 2 and 3 $+(a_{32}^{''})^{(6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6)}(T_{33},t)$ are sixth augmentation coefficient for category 1, 2 and 3 $+(a_{36}^{''})^{(7,7,7)}(T_{37},t)$, $+(a_{37}^{''})^{(7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7)}(T_{37},t)$ are seventh augmentation coefficient for category 1, 2 and 3 $+(a_{40}^{''})^{(8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8)}(T_{41},t)$ are eight augmentation coefficient for category 1, 2 and 3 $+(a_{44}^{''})^{(9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9)}(T_{45},t)$, $+(a_{46}^{''})^{(9,9)}(T_{45},t)$ are ninth augmentation coefficient for category 1, 2 and 3 | |
| $\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \begin{bmatrix} (b_{16}^{'})^{(2)}\,-(b_{16}^{''})^{(2)}(G_{19},t) & -(b_{13}^{''})^{(1,1,)}(G,t) & -(b_{20}^{''})^{(3,3,3,)}(G_{23},t) \\ -(b_{24}^{''})^{(4,4,4,4,4)}(G_{27},t) & -(b_{28}^{''})^{(5,5,5,5,5)}(G_{31},t) & -(b_{32}^{''})^{(6,6,6,6,6)}(G_{35},t) \\ -(b_{36}^{''})^{(7,7,7)}(G_{39},t) & -(b_{40}^{''})^{(8,8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9)}(G_{47},t) \end{bmatrix} T_{16}$ | *64* |
| $\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \begin{bmatrix} (b_{17}^{'})^{(2)}\,-(b_{17}^{''})^{(2)}(G_{19},t) & -(b_{14}^{''})^{(1,1,)}(G,t) & -(b_{21}^{''})^{(3,3,3,)}(G_{23},t) \\ -(b_{25}^{''})^{(4,4,4,4,4)}(G_{27},t) & -(b_{29}^{''})^{(5,5,5,5,5)}(G_{31},t) & -(b_{33}^{''})^{(6,6,6,6,6)}(G_{35},t) \\ -(b_{37}^{''})^{(7,7,7)}(G_{39},t) & -(b_{41}^{''})^{(8,8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9)}(G_{47},t) \end{bmatrix} T_{17}$ | *65* |
| $\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \begin{bmatrix} (b_{18}^{'})^{(2)}\,-(b_{18}^{''})^{(2)}(G_{19},t) & -(b_{15}^{''})^{(1,1,)}(G,t) & -(b_{22}^{''})^{(3,3,3,)}(G_{23},t) \\ -(b_{26}^{''})^{(4,4,4,4,4)}(G_{27},t) & -(b_{30}^{''})^{(5,5,5,5,5)}(G_{31},t) & -(b_{34}^{''})^{(6,6,6,6,6)}(G_{35},t) \\ -(b_{38}^{''})^{(7,7,7)}(G_{39},t) & -(b_{42}^{''})^{(8,8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9)}(G_{47},t) \end{bmatrix} T_{18}$ | *66* |
| where $-(b_{16}^{''})^{(2)}(G_{19},t)$, $-(b_{17}^{''})^{(2)}(G_{19},t)$, $-(b_{18}^{''})^{(2)}(G_{19},t)$ are first detrition coefficients for category 1, 2 and 3 $-(b_{13}^{''})^{(1,1,)}(G,t)$, $-(b_{14}^{''})^{(1,1,)}(G,t)$, $-(b_{15}^{''})^{(1,1,)}(G,t)$ are second detrition coefficients for category 1,2 and 3 $-(b_{20}^{''})^{(3,3,3,)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,)}(G_{23},t)$ are third detrition coefficients for category 1,2 and 3 $-(b_{24}^{''})^{(4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1,2 and 3 $-(b_{28}^{''})^{(5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1,2 and 3 $-(b_{32}^{''})^{(6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6)}(G_{35},t)$ are sixth detrition coefficients for category 1,2 and 3 $-(b_{36}^{''})^{(7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1,2 and 3 | |

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| $-(b_{40}^{''})^{(8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8)}(G_{43},t)$ are eight detrition coefficients for category 1,2 and 3<br>$-(b_{44}^{''})^{(9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1,2 and 3 | |
| $\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \left[\begin{array}{c} (a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21},t)\ +(a_{16}^{''})^{(2,2,2)}(T_{17},t)\ +(a_{13}^{''})^{(1,1,1,)}(T_{14},t) \\ +(a_{24}^{''})^{(4,4,4,4,4,4)}(T_{25},t)\ +(a_{28}^{''})^{(5,5,5,5,5,5)}(T_{29},t)\ +(a_{32}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ +(a_{36}^{''})^{(7,7,7,7)}(T_{37},t)\ +(a_{40}^{''})^{(8,8,8,8)}(T_{41},t)\ +(a_{44}^{''})^{(9,9,9)}(T_{45},t) \end{array}\right] G_{20}$ | 67 |
| $\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \left[\begin{array}{c} (a_{21}^{'})^{(3)}+(a_{21}^{''})^{(3)}(T_{21},t)\ +(a_{17}^{''})^{(2,2,2)}(T_{17},t)\ +(a_{14}^{''})^{(1,1,1,)}(T_{14},t) \\ +(a_{25}^{''})^{(4,4,4,4,4,4)}(T_{25},t)\ +(a_{29}^{''})^{(5,5,5,5,5,5)}(T_{29},t)\ +(a_{33}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ +(a_{37}^{''})^{(7,7,7,7)}(T_{37},t)\ +(a_{41}^{''})^{(8,8,8,8)}(T_{41},t)\ +(a_{45}^{''})^{(9,9,9)}(T_{45},t) \end{array}\right] G_{21}$ | 68 |
| $\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \left[\begin{array}{c} (a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21},t)\ +(a_{18}^{''})^{(2,2,2)}(T_{17},t)\ +(a_{15}^{''})^{(1,1,1,)}(T_{14},t) \\ +(a_{26}^{''})^{(4,4,4,4,4,4)}(T_{25},t)\ +(a_{30}^{''})^{(5,5,5,5,5,5)}(T_{29},t)\ +(a_{34}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ +(a_{38}^{''})^{(7,7,7,7)}(T_{37},t)\ +(a_{42}^{''})^{(8,8,8,8)}(T_{41},t)\ +(a_{46}^{''})^{(9,9,9)}(T_{45},t) \end{array}\right] G_{22}$ | 69 |
| $+(a_{20}^{''})^{(3)}(T_{21},t)$, $+(a_{21}^{''})^{(3)}(T_{21},t)$, $+(a_{22}^{''})^{(3)}(T_{21},t)$ are first augmentation coefficients for category 1, 2 and 3<br>$+(a_{16}^{''})^{(2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2)}(T_{17},t)$ are second augmentation coefficients for category 1, 2 and 3<br>$+(a_{13}^{''})^{(1,1,1,)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,)}(T_{14},t)$ are third augmentation coefficients for category 1, 2 and 3<br>$+(a_{24}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficients for category 1, 2 and 3<br>$+(a_{28}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficients for category 1, 2 and 3<br>$+(a_{32}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation coefficients for category 1, 2 and 3<br>$+(a_{36}^{''})^{(7,7,7,7)}(T_{37},t)$, $+(a_{37}^{''})^{(7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7)}(T_{37},t)$ are seventh augmentation coefficients for category 1, 2 and 3<br>$+(a_{40}^{''})^{(8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8)}(T_{41},t)$ are eight augmentation coefficients for category 1, 2 and 3<br>$+(a_{44}^{''})^{(9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9)}(T_{45},t)$, $+(a_{46}^{''})^{(9,9,9)}(T_{45},t)$ are ninth augmentation coefficients for category 1, 2 and 3 | |
| $\frac{dT_{20}}{dt} = (b_{20})^{(3)}T_{21} - \left[\begin{array}{c} (b_{20}^{'})^{(3)}-(b_{20}^{''})^{(3)}(G_{23},t)\ -(b_{16}^{''})^{(2,2,2)}(G_{19},t)\ -(b_{13}^{''})^{(1,1,1,)}(G,t) \\ -(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t)\ -(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t)\ -(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t) \\ -(b_{36}^{''})^{(7,7,7,7)}(G_{39},t)\ -(b_{40}^{''})^{(8,8,8,8)}(G_{43},t)\ -(b_{44}^{''})^{(9,9,9)}(G_{47},t) \end{array}\right] T_{20}$ | 70 |

$$\frac{dT_{21}}{dt} = (b_{21})^{(3)}T_{20} - \begin{bmatrix} (b'_{21})^{(3)} & -(b''_{21})^{(3)}(G_{23},t) & -(b''_{17})^{(2,2,2)}(G_{19},t) & -(b''_{14})^{(1,1,1,)}(G,t) \\ -(b''_{25})^{(4,4,4,4,4,4)}(G_{27},t) & -(b''_{29})^{(5,5,5,5,5,5)}(G_{31},t) & -(b''_{33})^{(6,6,6,6,6,6)}(G_{35},t) \\ -(b''_{37})^{(7,7,7,7)}(G_{39},t) & -(b''_{41})^{(8,8,8,8)}(G_{43},t) & -(b''_{45})^{(9,9,9)}(G_{47},t) \end{bmatrix} T_{21} \quad (71)$$

$$\frac{dT_{22}}{dt} = (b_{22})^{(3)}T_{21} - \begin{bmatrix} (b'_{22})^{(3)} & -(b''_{22})^{(3)}(G_{23},t) & -(b''_{18})^{(2,2,2)}(G_{19},t) & -(b''_{15})^{(1,1,1,)}(G,t) \\ -(b''_{26})^{(4,4,4,4,4,4)}(G_{27},t) & -(b''_{30})^{(5,5,5,5,5,5)}(G_{31},t) & -(b''_{34})^{(6,6,6,6,6,6)}(G_{35},t) \\ -(b''_{38})^{(7,7,7,7)}(G_{39},t) & -(b''_{42})^{(8,8,8,8)}(G_{43},t) & -(b''_{46})^{(9,9,9)}(G_{47},t) \end{bmatrix} T_{22} \quad (72)$$

$-(b''_{20})^{(3)}(G_{23},t)$, $-(b''_{21})^{(3)}(G_{23},t)$, $-(b''_{22})^{(3)}(G_{23},t)$ *are first detrition coefficients for category 1, 2 and 3*

$-(b''_{16})^{(2,2,2)}(G_{19},t)$, $-(b''_{17})^{(2,2,2)}(G_{19},t)$, $-(b''_{18})^{(2,2,2)}(G_{19},t)$ *are second detrition coefficients for category 1, 2 and 3*

$-(b''_{13})^{(1,1,1,)}(G,t)$, $-(b''_{14})^{(1,1,1,)}(G,t)$, $-(b''_{15})^{(1,1,1,)}(G,t)$ *are third detrition coefficients for category 1,2 and 3*

$-(b''_{24})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b''_{25})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b''_{26})^{(4,4,4,4,4,4)}(G_{27},t)$ *are fourth detrition coefficients for category 1, 2 and 3*

$-(b''_{28})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b''_{29})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b''_{30})^{(5,5,5,5,5,5)}(G_{31},t)$ *are fifth detrition coefficients for category 1, 2 and 3*

$-(b''_{32})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b''_{33})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b''_{34})^{(6,6,6,6,6,6)}(G_{35},t)$ *are sixth detrition coefficients for category 1, 2 and 3*

$-(b''_{36})^{(7,7,7,7)}(G_{39},t)$, $-(b''_{37})^{(7,7,7,7)}(G_{39},t)$ $-(b''_{38})^{(7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3

$-(b''_{40})^{(8,8,8,8)}(G_{43},t)$, $-(b''_{41})^{(8,8,8,8)}(G_{43},t)$, $-(b''_{42})^{(8,8,8,8)}(G_{43},t)$ are eight detrition coefficients for category 1, 2 and 3

$-(b''_{46})^{(9,9,9)}(G_{47},t)$, $-(b''_{45})^{(9,9,9)}(G_{47},t)$, $-(b''_{44})^{(9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3

$$\frac{dG_{24}}{dt} = (a_{24})^{(4)}G_{25} - \begin{bmatrix} (a'_{24})^{(4)} & +(a''_{24})^{(4)}(T_{25},t) & +(a''_{28})^{(5,5,)}(T_{29},t) & +(a''_{32})^{(6,6,)}(T_{33},t) \\ +(a''_{13})^{(1,1,1,1)}(T_{14},t) & +(a''_{16})^{(2,2,2,2)}(T_{17},t) & +(a''_{20})^{(3,3,3,3)}(T_{21},t) \\ +(a''_{36})^{(7,7,7,7,7)}(T_{37},t) & +(a''_{40})^{(8,8,8,8,8)}(T_{41},t) & +(a''_{44})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{24} \quad (73)$$

$$\frac{dG_{25}}{dt} = (a_{25})^{(4)}G_{24} - \begin{bmatrix} (a'_{25})^{(4)} & +(a''_{25})^{(4)}(T_{25},t) & +(a''_{29})^{(5,5,)}(T_{29},t) & +(a''_{33})^{(6,6)}(T_{33},t) \\ +(a''_{14})^{(1,1,1,1)}(T_{14},t) & +(a''_{17})^{(2,2,2,2)}(T_{17},t) & +(a''_{21})^{(3,3,3,3)}(T_{21},t) \\ +(a''_{37})^{(7,7,7,7,7)}(T_{37},t) & +(a''_{41})^{(8,8,8,8,8)}(T_{41},t) & +(a''_{45})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{25} \quad (74)$$

$$\frac{dG_{26}}{dt} = (a_{26})^{(4)}G_{25} - \begin{bmatrix} (a'_{26})^{(4)} & +(a''_{26})^{(4)}(T_{25},t) & +(a''_{30})^{(5,5,)}(T_{29},t) & +(a''_{34})^{(6,6,)}(T_{33},t) \\ +(a''_{15})^{(1,1,1,1)}(T_{14},t) & +(a''_{18})^{(2,2,2,2)}(T_{17},t) & +(a''_{22})^{(3,3,3,3)}(T_{21},t) \\ +(a''_{38})^{(7,7,7,7,7)}(T_{37},t) & +(a''_{42})^{(8,8,8,8,8)}(T_{41},t) & +(a''_{46})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{26} \quad (75)$$

$(a''_{24})^{(4)}(T_{25},t)$, $(a''_{25})^{(4)}(T_{25},t)$, $(a''_{26})^{(4)}(T_{25},t)$ *are first augmentation coefficients category* 1, 2 3

$+(a_{28}^{''})^{(5,5,)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,)}(T_{29},t)$ *are second augmentation coefficient for category* $1,2$ *and* $3$

$+(a_{32}^{''})^{(6,6,)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,)}(T_{33},t)$ *are third augmentation coefficient for category* $1,2$ *and* $3$

$+(a_{13}^{''})^{(1,1,1,1)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1)}(T_{14},t)$ *are fourth augmentation coefficients for category* $1,2$ *and* $3$

$+(a_{16}^{''})^{(2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2)}(T_{17},t)$ *are fifth augmentation coefficients for category* $1,2$ *and* $3$

$+(a_{20}^{''})^{(3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3)}(T_{21},t)$ *are sixth augmentation coefficients for category* $1,2$ *and* $3$

$+(a_{36}^{''})^{(7,7,7,7,7)}(T_{37},t)$, $+(a_{37}^{''})^{(7,7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7,7)}(T_{37},t)$ *are seventh augmentation coefficients for category* $1,2$ *and* $3$

$+(a_{40}^{''})^{(8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8,8)}(T_{41},t)$ *are eighth augmentation coefficients for category* $1,2$ *and* $3$

$+(a_{46}^{''})^{(9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9)}(T_{45},t)$ are ninth detrition coefficients for category 1 2 3

$$\frac{dT_{24}}{dt}=(b_{24})^{(4)}T_{25}-\begin{bmatrix}(b_{24}^{'})^{(4)}&-(b_{24}^{''})^{(4)}(G_{27},t)&-(b_{28}^{''})^{(5,5,)}(G_{31},t)&-(b_{32}^{''})^{(6,6,)}(G_{35},t)\\-(b_{13}^{''})^{(1,1,1,1)}(G,t)&-(b_{16}^{''})^{(2,2,2,2)}(G_{19},t)&-(b_{20}^{''})^{(3,3,3,3)}(G_{23},t)\\-(b_{36}^{''})^{(7,7,7,7,7)}(G_{39},t)&-(b_{40}^{''})^{(8,8,8,8,8)}(G_{43},t)&-(b_{44}^{''})^{(9,9,9,9)}(G_{47},t)\end{bmatrix}T_{24} \quad (76)$$

$$\frac{dT_{25}}{dt}=(b_{25})^{(4)}T_{24}-\begin{bmatrix}(b_{25}^{'})^{(4)}&-(b_{25}^{''})^{(4)}(G_{27},t)&-(b_{29}^{''})^{(5,5,)}(G_{31},t)&-(b_{33}^{''})^{(6,6,)}(G_{35},t)\\-(b_{14}^{''})^{(1,1,1,1)}(G,t)&-(b_{17}^{''})^{(2,2,2,2)}(G_{19},t)&-(b_{21}^{''})^{(3,3,3,3)}(G_{23},t)\\-(b_{37}^{''})^{(7,7,7,7,7)}(G_{39},t)&-(b_{41}^{''})^{(8,8,8,8,8)}(G_{43},t)&-(b_{45}^{''})^{(9,9,9,9)}(G_{47},t)\end{bmatrix}T_{25} \quad (77)$$

$$\frac{dT_{26}}{dt}=(b_{26})^{(4)}T_{25}-\begin{bmatrix}(b_{26}^{'})^{(4)}&-(b_{26}^{''})^{(4)}(G_{27},t)&-(b_{30}^{''})^{(5,5,)}(G_{31},t)&-(b_{34}^{''})^{(6,6,)}(G_{35},t)\\-(b_{15}^{''})^{(1,1,1,1)}(G,t)&-(b_{18}^{''})^{(2,2,2,2)}(G_{19},t)&-(b_{22}^{''})^{(3,3,3,3)}(G_{23},t)\\-(b_{38}^{''})^{(7,7,7,7,7)}(G_{39},t)&-(b_{42}^{''})^{(8,8,8,8,8)}(G_{43},t)&-(b_{46}^{''})^{(9,9,9,9)}(G_{47},t)\end{bmatrix}T_{26} \quad (78)$$

*Where* $-(b_{24}^{''})^{(4)}(G_{27},t)$, $-(b_{25}^{''})^{(4)}(G_{27},t)$, $-(b_{26}^{''})^{(4)}(G_{27},t)$ *are first detrition coefficients for category* $1,2$ *and* $3$

$-(b_{28}^{''})^{(5,5,)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,)}(G_{31},t)$ *are second detrition coefficients for category* $1,2$ *and* $3$

$-(b_{32}^{''})^{(6,6,)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,)}(G_{35},t)$ *are third detrition coefficients for category* $1,2$ *and* $3$

$-(b_{13}^{''})^{(1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1)}(G,t)$, $-(b_{15}^{''})^{(1,1,1,1)}(G,t)$ *are fourth detrition coefficients for category* $1,2$ *and* $3$

$-(b_{16}^{''})^{(2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2)}(G_{19},t)$ *are fifth detrition coefficients for category* $1,2$ *and* $3$

$-(b_{20}^{''})^{(3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3)}(G_{23},t)$

| | |
|---|---|
| *are sixth detrition coefficients for category* 1, 2 *and* 3 $\boxed{-(b_{36}^{''})^{(7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{37}^{''})^{(7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{38}^{''})^{(7,7,7,7,7)}(G_{39},t)}$ *are seventh detrition coefficients for category* 1, 2 *and* 3 $\boxed{-(b_{40}^{''})^{(8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{41}^{''})^{(8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{42}^{''})^{(8,8,8,8,8)}(G_{43},t)}$ *are eighth detrition coefficients for category* 1, 2 *and* 3 $\boxed{-(b_{46}^{''})^{(9,9,9,9)}(G_{47},t)}$, $\boxed{-(b_{45}^{''})^{(9,9,9,9)}(G_{47},t)}$, $\boxed{-(b_{44}^{''})^{(9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1 2 3 | |
| $\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \left[\begin{array}{ccc} (a_{28}^{'})^{(5)}\boxed{+(a_{28}^{''})^{(5)}(T_{29},t)} & \boxed{+(a_{24}^{''})^{(4,4,)}(T_{25},t)} & \boxed{+(a_{32}^{''})^{(6,6,6)}(T_{33},t)} \\ \boxed{+(a_{13}^{''})^{(1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{16}^{''})^{(2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{20}^{''})^{(3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{36}^{''})^{(7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{40}^{''})^{(8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{44}^{''})^{(9,9,9,9,9)}(T_{45},t)} \end{array}\right] G_{28}$ | *79* |
| $\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \left[\begin{array}{ccc} (a_{29}^{'})^{(5)}\boxed{+(a_{29}^{''})^{(5)}(T_{29},t)} & \boxed{+(a_{25}^{''})^{(4,4,)}(T_{25},t)} & \boxed{+(a_{33}^{''})^{(6,6,6)}(T_{33},t)} \\ \boxed{+(a_{14}^{''})^{(1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{17}^{''})^{(2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{21}^{''})^{(3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{37}^{''})^{(7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{41}^{''})^{(8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{45}^{''})^{(9,9,9,9,9)}(T_{45},t)} \end{array}\right] G_{29}$ | *80* |
| $\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \left[\begin{array}{ccc} (a_{30}^{'})^{(5)}\boxed{+(a_{30}^{''})^{(5)}(T_{29},t)} & \boxed{+(a_{26}^{''})^{(4,4,)}(T_{25},t)} & \boxed{+(a_{34}^{''})^{(6,6,6)}(T_{33},t)} \\ \boxed{+(a_{15}^{''})^{(1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{18}^{''})^{(2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{22}^{''})^{(3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{38}^{''})^{(7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{42}^{''})^{(8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{46}^{''})^{(9,9,9,9,9)}(T_{45},t)} \end{array}\right] G_{30}$ | *81* |
| *Where* $\boxed{+(a_{28}^{''})^{(5)}(T_{29},t)}$, $\boxed{+(a_{29}^{''})^{(5)}(T_{29},t)}$, $\boxed{+(a_{30}^{''})^{(5)}(T_{29},t)}$ *are first augmentation coefficients for category* 1, 2 *and* 3 *And* $\boxed{+(a_{24}^{''})^{(4,4,)}(T_{25},t)}$, $\boxed{+(a_{25}^{''})^{(4,4,)}(T_{25},t)}$, $\boxed{+(a_{26}^{''})^{(4,4,)}(T_{25},t)}$ *are second augmentation coefficient for category* 1, 2 *and* 3 $\boxed{+(a_{32}^{''})^{(6,6,6)}(T_{33},t)}$, $\boxed{+(a_{33}^{''})^{(6,6,6)}(T_{33},t)}$, $\boxed{+(a_{34}^{''})^{(6,6,6)}(T_{33},t)}$ *are third augmentation coefficient for category* 1, 2 *and* 3 $\boxed{+(a_{13}^{''})^{(1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a_{14}^{''})^{(1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a_{15}^{''})^{(1,1,1,1,1)}(T_{14},t)}$ are fourth augmentation coefficients for category 1,2, and 3 $\boxed{+(a_{16}^{''})^{(2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a_{17}^{''})^{(2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a_{18}^{''})^{(2,2,2,2,2)}(T_{17},t)}$ are fifth augmentation coefficients for category 1,2,and 3 $\boxed{+(a_{20}^{''})^{(3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a_{21}^{''})^{(3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a_{22}^{''})^{(3,3,3,3,3)}(T_{21},t)}$ are sixth augmentation coefficients for category 1,2, 3 $\boxed{+(a_{36}^{''})^{(7,7,7,7,7,7)}(T_{37},t)}$, $\boxed{+(a_{37}^{''})^{(7,7,7,7,7,7)}(T_{37},t)}$, $\boxed{+(a_{38}^{''})^{(7,7,7,7,7,7)}(T_{37},t)}$ are seventh augmentation coefficients for category 1,2, 3 $\boxed{+(a_{40}^{''})^{(8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{41}^{''})^{(8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{42}^{''})^{(8,8,8,8,8,8)}(T_{41},t)}$ are eighth augmentation coefficients for category 1,2, 3 $\boxed{+(a_{46}^{''})^{(9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{45}^{''})^{(9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{44}^{''})^{(9,9,9,9,9)}(T_{45},t)}$ are ninth augmentation coefficients for category 1,2, 3 | |

$$\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \begin{bmatrix} (b'_{28})^{(5)} & -(b''_{28})^{(5)}(G_{31},t) & -(b''_{24})^{(4,4,)}(G_{27},t) & -(b''_{32})^{(6,6,6)}(G_{35},t) \\ -(b''_{13})^{(1,1,1,1,1)}(G,t) & -(b''_{16})^{(2,2,2,2,2)}(G_{19},t) & -(b''_{20})^{(3,3,3,3,3)}(G_{23},t) & \\ -(b''_{36})^{(7,7,7,7,7,7)}(G_{39},t) & -(b''_{40})^{(8,8,8,8,8,8)}(G_{43},t) & -(b''_{44})^{(9,9,9,9,9)}(G_{47},t) & \end{bmatrix} T_{28} \quad (82)$$

$$\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \begin{bmatrix} (b'_{29})^{(5)} & -(b''_{29})^{(5)}(G_{31},t) & -(b''_{25})^{(4,4,)}(G_{27},t) & -(b''_{33})^{(6,6,6)}(G_{35},t) \\ -(b''_{14})^{(1,1,1,1,1)}(G,t) & -(b''_{17})^{(2,2,2,2,2)}(G_{19},t) & -(b''_{21})^{(3,3,3,3,3)}(G_{23},t) & \\ -(b''_{37})^{(7,7,7,7,7,7)}(G_{39},t) & -(b''_{41})^{(8,8,8,8,8,8)}(G_{43},t) & -(b''_{45})^{(9,9,9,9,9)}(G_{47},t) & \end{bmatrix} T_{29} \quad (83)$$

$$\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \begin{bmatrix} (b'_{30})^{(5)} & -(b''_{30})^{(5)}(G_{31},t) & -(b''_{26})^{(4,4,)}(G_{27},t) & -(b''_{34})^{(6,6,6)}(G_{35},t) \\ -(b''_{15})^{(1,1,1,1,1,)}(G,t) & -(b''_{18})^{(2,2,2,2,2)}(G_{19},t) & -(b''_{22})^{(3,3,3,3,3)}(G_{23},t) & \\ -(b''_{38})^{(7,7,7,7,7,7)}(G_{39},t) & -(b''_{42})^{(8,8,8,8,8,8)}(G_{43},t) & -(b''_{46})^{(9,9,9,9,9)}(G_{47},t) & \end{bmatrix} T_{30} \quad (84)$$

*where* $-(b''_{28})^{(5)}(G_{31},t)$ , $-(b''_{29})^{(5)}(G_{31},t)$ , $-(b''_{30})^{(5)}(G_{31},t)$ *are first detrition coefficients for category* 1, 2 *and* 3

$-(b''_{24})^{(4,4,)}(G_{27},t)$, $-(b''_{25})^{(4,4,)}(G_{27},t)$, $-(b''_{26})^{(4,4,)}(G_{27},t)$ *are second detrition coefficients for category* 1,2 *and* 3

$-(b''_{32})^{(6,6,6)}(G_{35},t)$, $-(b''_{33})^{(6,6,6)}(G_{35},t)$, $-(b''_{34})^{(6,6,6)}(G_{35},t)$ *are third detrition coefficients for category* 1,2 *and* 3

$-(b''_{13})^{(1,1,1,1,1)}(G,t)$, $-(b''_{14})^{(1,1,1,1,1)}(G,t)$ , $-(b''_{15})^{(1,1,1,1,1,)}(G,t)$ are fourth detrition coefficients for category 1,2, and 3

$-(b''_{16})^{(2,2,2,2,2)}(G_{19},t)$, $-(b''_{17})^{(2,2,2,2,2)}(G_{19},t)$, $-(b''_{18})^{(2,2,2,2,2)}(G_{19},t)$ are fifth detrition coefficients for category 1,2, and 3

$-(b''_{20})^{(3,3,3,3,3)}(G_{23},t)$, $-(b''_{21})^{(3,3,3,3,3)}(G_{23},t)$, $-(b''_{22})^{(3,3,3,3,3)}(G_{23},t)$ are sixth detrition coefficients for category 1,2, and 3

$-(b''_{36})^{(7,7,7,7,7,7)}(G_{39},t)$, $-(b''_{37})^{(7,7,7,7,7,7)}(G_{39},t)$, $-(b''_{38})^{(7,7,7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1,2, and 3

$-(b''_{42})^{(8,8,8,8,8,8)}(G_{43},t)$, $-(b''_{41})^{(8,8,8,8,8,8)}(G_{43},t)$, $-(b''_{40})^{(8,8,8,8,8,8)}(G_{43},t)$ are eighth detrition coefficients for category 1,2, and 3

$-(b''_{46})^{(9,9,9,9,9)}(G_{47},t)$, $-(b''_{45})^{(9,9,9,9,9)}(G_{47},t)$, $-(b''_{44})^{(9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1,2, and 3

$$\frac{dG_{32}}{dt} = (a_{32})^{(6)}G_{33} - \begin{bmatrix} (a'_{32})^{(6)} & +(a''_{32})^{(6)}(T_{33},t) & +(a''_{28})^{(5,5,5)}(T_{29},t) & +(a''_{24})^{(4,4,4,)}(T_{25},t) \\ +(a''_{13})^{(1,1,1,1,1,1)}(T_{14},t) & +(a''_{16})^{(2,2,2,2,2,2)}(T_{17},t) & +(a''_{20})^{(3,3,3,3,3,3)}(T_{21},t) & \\ +(a''_{36})^{(7,7,7,7,7,7,7)}(T_{37},t) & +(a''_{40})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a''_{44})^{(9,9,9,9,9,9)}(T_{45},t) & \end{bmatrix} G_{32} \quad (85)$$

$$\frac{dG_{33}}{dt} = (a_{33})^{(6)}G_{32} - \begin{bmatrix} (a'_{33})^{(6)} & +(a''_{33})^{(6)}(T_{33},t) & +(a''_{29})^{(5,5,5)}(T_{29},t) & +(a''_{25})^{(4,4,4,)}(T_{25},t) \\ +(a''_{14})^{(1,1,1,1,1,1)}(T_{14},t) & +(a''_{17})^{(2,2,2,2,2,2)}(T_{17},t) & +(a''_{21})^{(3,3,3,3,3,3)}(T_{21},t) & \\ +(a''_{37})^{(7,7,7,7,7,7,7)}(T_{37},t) & +(a''_{41})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a''_{45})^{(9,9,9,9,9,9)}(T_{45},t) & \end{bmatrix} G_{33} \quad (86)$$

$$\frac{dG_{34}}{dt} = (a_{34})^{(6)} G_{33} - \begin{bmatrix} (a'_{34})^{(6)} & +(a''_{34})^{(6)}(T_{33},t) & +(a''_{30})^{(5,5,5)}(T_{29},t) & +(a''_{26})^{(4,4,4,)}(T_{25},t) \\ & +(a''_{15})^{(1,1,1,1,1,1)}(T_{14},t) & +(a''_{18})^{(2,2,2,2,2,2)}(T_{17},t) & +(a''_{22})^{(3,3,3,3,3,3)}(T_{21},t) \\ & +(a''_{38})^{(7,7,7,7,7,7,7)}(T_{37},t) & +(a''_{42})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a''_{46})^{(9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{34} \quad (87)$$

$+(a''_{32})^{(6)}(T_{33},t)$, $+(a''_{33})^{(6)}(T_{33},t)$, $+(a''_{34})^{(6)}(T_{33},t)$ *are first augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{28})^{(5,5,5)}(T_{29},t)$, $+(a''_{29})^{(5,5,5)}(T_{29},t)$, $+(a''_{30})^{(5,5,5)}(T_{29},t)$ *are second augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{24})^{(4,4,4,)}(T_{25},t)$, $+(a''_{25})^{(4,4,4,)}(T_{25},t)$, $+(a''_{26})^{(4,4,4,)}(T_{25},t)$ *are third augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{13})^{(1,1,1,1,1,1)}(T_{14},t)$, $+(a''_{14})^{(1,1,1,1,1,1)}(T_{14},t)$, $+(a''_{15})^{(1,1,1,1,1,1)}(T_{14},t)$ - are fourth augmentation coefficients

$+(a''_{16})^{(2,2,2,2,2,2)}(T_{17},t)$, $+(a''_{17})^{(2,2,2,2,2,2)}(T_{17},t)$, $+(a''_{18})^{(2,2,2,2,2,2)}(T_{17},t)$ - fifth augmentation coefficients

$+(a''_{20})^{(3,3,3,3,3,3)}(T_{21},t)$, $+(a''_{21})^{(3,3,3,3,3,3)}(T_{21},t)$, $+(a''_{22})^{(3,3,3,3,3,3)}(T_{21},t)$ sixth augmentation coefficients

$+(a''_{36})^{(7,7,7,7,7,7,7)}(T_{37},t)$, $+(a''_{37})^{(7,7,7,7,7,7,7)}(T_{37},t)$, $+(a''_{38})^{(7,7,7,7,7,7,7)}(T_{37},t)$ seventh augmentation coefficients

$+(a''_{40})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a''_{41})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a''_{42})^{(8,8,8,8,8,8,8)}(T_{41},t)$ Eighth augmentation coefficients

$+(a''_{44})^{(9,9,9,9,9,9)}(T_{45},t)$, $+(a''_{45})^{(9,9,9,9,9,9)}(T_{45},t)$, $+(a''_{46})^{(9,9,9,9,9,9)}(T_{45},t)$ ninth augmentation coefficients

$$\frac{dT_{32}}{dt} = (b_{32})^{(6)} T_{33} - \begin{bmatrix} (b'_{32})^{(6)} & -(b''_{32})^{(6)}(G_{35},t) & -(b''_{28})^{(5,5,5)}(G_{31},t) & -(b''_{24})^{(4,4,4,)}(G_{27},t) \\ & -(b''_{13})^{(1,1,1,1,1,1)}(G,t) & -(b''_{16})^{(2,2,2,2,2,2)}(G_{19},t) & -(b''_{20})^{(3,3,3,3,3,3)}(G_{23},t) \\ & -(b''_{36})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b''_{40})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b''_{44})^{(9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{32} \quad (88)$$

$$\frac{dT_{33}}{dt} = (b_{33})^{(6)} T_{32} - \begin{bmatrix} (b'_{33})^{(6)} & -(b''_{33})^{(6)}(G_{35},t) & -(b''_{29})^{(5,5,5)}(G_{31},t) & -(b''_{25})^{(4,4,4,)}(G_{27},t) \\ & -(b''_{14})^{(1,1,1,1,1,1)}(G,t) & -(b''_{17})^{(2,2,2,2,2,2)}(G_{19},t) & -(b''_{21})^{(3,3,3,3,3,3)}(G_{23},t) \\ & -(b''_{37})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b''_{41})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b''_{45})^{(9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{33} \quad (89)$$

$$\frac{dT_{34}}{dt} = (b_{34})^{(6)} T_{33} - \begin{bmatrix} (b'_{34})^{(6)} & -(b''_{34})^{(6)}(G_{35},t) & -(b''_{30})^{(5,5,5)}(G_{31},t) & -(b''_{26})^{(4,4,4,)}(G_{27},t) \\ & -(b''_{15})^{(1,1,1,1,1,1)}(G,t) & -(b''_{18})^{(2,2,2,2,2,2)}(G_{19},t) & -(b''_{22})^{(3,3,3,3,3,3)}(G_{23},t) \\ & -(b''_{38})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b''_{42})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b''_{46})^{(9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{34} \quad (90)$$

$-(b''_{32})^{(6)}(G_{35},t)$, $-(b''_{33})^{(6)}(G_{35},t)$, $-(b''_{34})^{(6)}(G_{35},t)$ *are first detrition coefficients for category* 1, 2 *and* 3

$-(b''_{28})^{(5,5,5)}(G_{31},t)$, $-(b''_{29})^{(5,5,5)}(G_{31},t)$, $-(b''_{30})^{(5,5,5)}(G_{31},t)$ *are second detrition coefficients for category* 1, 2 *and* 3

$-(b_{24}^{''})^{(4,4,4,)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,)}(G_{27},t)$ *are third detrition coefficients for category* 1,2 *and* 3

$-(b_{13}^{''})^{(1,1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1,1)}(G,t)$, $-(b_{15}^{''})^{(1,1,1,1,1,1)}(G,t)$ are fourth detrition coefficients for category 1, 2, and 3

$-(b_{16}^{''})^{(2,2,2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2,2,2)}(G_{19},t)$ are fifth detrition coefficients for category 1, 2, and 3

$-(b_{20}^{''})^{(3,3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3,3)}(G_{23},t)$ are sixth detrition coefficients for category 1, 2, and 3

$-(b_{36}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2, and 3

$-(b_{40}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t)$ are eighth detrition coefficients for category 1, 2, and 3

$-(b_{46}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2, and 3

$$\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \begin{bmatrix} (a_{36}^{'})^{(7)}+(a_{36}^{''})^{(7)}(T_{37},t) & +(a_{16}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t) & +(a_{20}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{24}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t) & +(a_{28}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t) & +(a_{32}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{13}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t) & +(a_{40}^{''})^{(8,8,8,8,8,8,8,8,)}(T_{41},t) & +(a_{44}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{13} \quad (91)$$

$$\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \begin{bmatrix} (a_{37}^{'})^{(7)}+(a_{37}^{''})^{(7)}(T_{37},t) & +(a_{17}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t) & +(a_{21}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{25}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t) & +(a_{29}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t) & +(a_{33}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{13}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t) & +(a_{41}^{''})^{(8,8,8,8,8,8,8,8)}(T_{41},t) & +(a_{45}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{14} \quad (92)$$

$$\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \begin{bmatrix} (a_{38}^{'})^{(7)}+(a_{38}^{''})^{(7)}(T_{37},t) & +(a_{18}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t) & +(a_{22}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{26}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t) & +(a_{30}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t) & +(a_{34}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{15}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t) & +(a_{42}^{''})^{(8,8,8,8,8,8,8,8)}(T_{41},t) & +(a_{46}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{15} \quad (93)$$

Where $(a_{36}^{''})^{(7)}(T_{37},t)$, $(a_{37}^{''})^{(7)}(T_{37},t)$, $(a_{38}^{''})^{(7)}(T_{37},t)$ are first augmentation coefficients for category 1, 2 and 3

$+(a_{16}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t)$ are second augmentation coefficient for category 1, 2 and 3

$+(a_{20}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3

$+(a_{24}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3

$+(a_{28}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation

coefficient for category 1, 2 and 3 $+(a_{32}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation coefficient for category 1, 2 and 3 $+(a_{13}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{13}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t)$ are seventh augmentation coefficient for category 1, 2 and 3 $+(a_{42}^{''})^{(8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{40}^{''})^{(8,8,8,8,8,8,8,8,)}(T_{41},t)$ are eighth augmentation coefficient for 1,2,3 $+(a_{46}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{36}}{dt} =$$

$$(b_{36})^{(7)}T_{37} - \begin{bmatrix} (b_{36}^{'})^{(7)}\left[-(b_{36}^{''})^{(7)}(G_{39},t)\right] & \left[-(b_{16}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)\right] & \left[-(b_{20}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)\right] \\ \left[-(b_{24}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t)\right] & \left[-(b_{28}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t)\right] & \left[-(b_{32}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t)\right] \\ \left[-(b_{13}^{''})^{(1,1,1,1,1,1,1)}(G,t)\right] & \left[-(b_{40}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t)\right] & \left[-(b_{44}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{13}$$

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$$\frac{dT_{37}}{dt} =$$

$$(b_{37})^{(7)}T_{36} - \begin{bmatrix} (b_{37}^{'})^{(7)}\left[-(b_{37}^{''})^{(7)}(G_{39},t)\right] & \left[-(b_{17}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)\right] & \left[-(b_{21}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)\right] \\ \left[-(b_{25}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t)\right] & \left[-(b_{29}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t)\right] & \left[-(b_{33}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t)\right] \\ \left[-(b_{14}^{''})^{(1,1,1,1,1,1,1)}(G,t)\right] & \left[-(b_{41}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t)\right] & \left[-(b_{45}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{14}$$

$$\frac{dT_{38}}{dt} =$$

$$(b_{38})^{(7)}T_{37} - \begin{bmatrix} (b_{38}^{'})^{(7)}\left[-(b_{38}^{''})^{(7)}(G_{39},t)\right] & \left[-(b_{18}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)\right] & \left[-(b_{22}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)\right] \\ \left[-(b_{26}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t)\right] & \left[-(b_{30}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t)\right] & \left[-(b_{34}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t)\right] \\ \left[-(b_{15}^{''})^{(1,1,1,1,1,1,1)}(G,t)\right] & \left[-(b_{42}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t)\right] & \left[-(b_{46}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{15}$$

Where $-(b_{36}^{''})^{(7)}(G_{39},t)$, $-(b_{37}^{''})^{(7)}(G_{39},t)$, $-(b_{38}^{''})^{(7)}(G_{39},t)$ are first detrition coefficients for category 1, 2 and 3 $-(b_{16}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3 $-(b_{20}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3 $-(b_{24}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3 $-(b_{28}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3 $-(b_{32}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t)$ are sixth detrition coefficients for category 1, 2 and 3 $-(b_{15}^{''})^{(1,1,1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1,1,1)}(G,t)$, $-(b_{13}^{''})^{(1,1,1,1,1,1,1)}(G,t)$ are seventh detrition coefficients for category 1, 2 and 3

$-(b_{40}'')^{(8,8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{41}'')^{(8,8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{42}'')^{(8,8,8,8,8,8,8,8)}(G_{43},t)$ are eighth detrition coefficients for category 1, 2 and 3

$-(b_{46}'')^{(9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}'')^{(9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{44}'')^{(9,9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3

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$$\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \begin{bmatrix} (a_{40}')^{(8)}+(a_{40}'')^{(8)}(T_{41},t) & +(a_{16}'')^{(2,2,2,2,2,2,2,2)}(T_{17},t) & +(a_{20}'')^{(3,3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{24}'')^{(4,4,4,4,4,4,4,4)}(T_{25},t) & +(a_{28}'')^{(5,5,5,5,5,5,5,5)}(T_{29},t) & +(a_{32}'')^{(6,6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{13}'')^{(1,1,1,1,1,1,1,1)}(T_{14},t) & +(a_{36}'')^{(7,7,7,7,7,7,7,7)}(T_{37},t) & +(a_{44}'')^{(9,9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{13}$$

$$\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \begin{bmatrix} (a_{41}')^{(8)}+(a_{41}'')^{(8)}(T_{41},t) & +(a_{17}'')^{(2,2,2,2,2,2,2,2)}(T_{17},t) & +(a_{21}'')^{(3,3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{25}'')^{(4,4,4,4,4,4,4,4)}(T_{25},t) & +(a_{29}'')^{(5,5,5,5,5,5,5,5)}(T_{29},t) & +(a_{33}'')^{(6,6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{13}'')^{(1,1,1,1,1,1,1,1)}(T_{14},t) & +(a_{37}'')^{(7,7,7,7,7,7,7,7)}(T_{37},t) & +(a_{45}'')^{(9,9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{14}$$

$$\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \begin{bmatrix} (a_{42}')^{(8)}+(a_{42}'')^{(8)}(T_{41},t) & +(a_{18}'')^{(2,2,2,2,2,2,2,2)}(T_{17},t) & +(a_{22}'')^{(3,3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{26}'')^{(4,4,4,4,4,4,4,4)}(T_{25},t) & +(a_{30}'')^{(5,5,5,5,5,5,5,5)}(T_{29},t) & +(a_{34}'')^{(6,6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{15}'')^{(1,1,1,1,1,1,1,1)}(T_{14},t) & +(a_{38}'')^{(7,7,7,7,7,7,7,7)}(T_{37},t) & +(a_{46}'')^{(9,9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{15}$$

Where $+(a_{40}'')^{(8)}(T_{41},t)$, $+(a_{41}'')^{(8)}(T_{41},t)$, $+(a_{42}'')^{(8)}(T_{41},t)$ are first augmentation coefficients for category 1, 2 and 3

$+(a_{16}'')^{(2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{17}'')^{(2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{18}'')^{(2,2,2,2,2,2,2,2)}(T_{17},t)$ are second augmentation coefficient for category 1, 2 and 3

$+(a_{20}'')^{(3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{21}'')^{(3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{22}'')^{(3,3,3,3,3,3,3,3)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3

$+(a_{24}'')^{(4,4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}'')^{(4,4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}'')^{(4,4,4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3

$+(a_{28}'')^{(5,5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}'')^{(5,5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}'')^{(5,5,5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficient for category 1, 2 and 3

$+(a_{32}'')^{(6,6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}'')^{(6,6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}'')^{(6,6,6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation coefficient for category 1, 2 and 3

$+(a_{13}'')^{(1,1,1,1,1,1,1,1)}(T_{14},t)$ $+(a_{14}'')^{(1,1,1,1,1,1,1,1)}(T_{14},t)$ $+(a_{15}'')^{(1,1,1,1,1,1,1,1)}(T_{14},t)$ are seventh augmentation coefficient for 1,2,3

$+(a_{36}'')^{(7,7,7,7,7,7,7,7)}(T_{37},t)$, $+(a_{37}'')^{(7,7,7,7,7,7,7,7)}(T_{37},t)$, $+(a_{38}'')^{(7,7,7,7,7,7,7,7)}(T_{37},t)$ are eighth augmentation coefficient for 1,2,3

$+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{40}}{dt} = (b_{40})^{(8)}T_{41} - \begin{bmatrix} (b_{40}^{'})^{(8)}\left[-(b_{40}^{''})^{(8)}(G_{43},t)\right] & -(b_{16}^{''})^{(2,2,2,2,2,2,2,2)}(G_{19},t) & -(b_{20}^{''})^{(3,3,3,3,3,3,3,3)}(G_{23},t) \\ -(b_{24}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t) & -(b_{28}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t) & -(b_{32}^{''})^{(6,6,6,6,6,6,6,6)}(G_{35},t) \\ -(b_{13}^{''})^{(1,1,1,1,1,1,1,1)}(G,t) & -(b_{36}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t) & -(b_{44}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{13}$$

$$\frac{dT_{41}}{dt} = (b_{41})^{(8)}T_{40} - \begin{bmatrix} (b_{41}^{'})^{(8)}\left[-(b_{41}^{''})^{(8)}(G_{43},t)\right] & -(b_{17}^{''})^{(2,2,2,2,2,2,2,2)}(G_{19},t) & -(b_{21}^{''})^{(3,3,3,3,3,3,3,3)}(G_{23},t) \\ -(b_{25}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t) & -(b_{29}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t) & -(b_{33}^{''})^{(6,6,6,6,6,6,6,6)}(G_{35},t) \\ -(b_{14}^{''})^{(1,1,1,1,1,1,1,1)}(G,t) & -(b_{37}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t) & -(b_{45}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{14}$$

$$\frac{dT_{42}}{dt} = (b_{42})^{(8)}T_{41} - \begin{bmatrix} (b_{42}^{'})^{(8)}\left[-(b_{42}^{''})^{(8)}(G_{43},t)\right] & -(b_{18}^{''})^{(2,2,2,2,2,2,2,2)}(G_{19},t) & -(b_{22}^{''})^{(3,3,3,3,3,3,3,3)}(G_{23},t) \\ -(b_{26}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t) & -(b_{30}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t) & -(b_{34}^{''})^{(6,6,6,6,6,6,6,6)}(G_{35},t) \\ -(b_{15}^{''})^{(1,1,1,1,1,1,1,1)}(G,t) & -(b_{38}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t) & -(b_{46}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{15}$$

Where $-(b_{36}^{''})^{(7)}(G_{39},t)$, $-(b_{37}^{''})^{(7)}(G_{39},t)$, $-(b_{38}^{''})^{(7)}(G_{39},t)$ are first detrition coefficients for category 1, 2 and 3

$-(b_{16}^{''})^{(2,2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2,2,2,2,2)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3

$-(b_{20}^{''})^{(3,3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3,3,3,3)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3

$-(b_{24}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3

$-(b_{28}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3

$-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{15}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$ are sixth detrition coefficients for category 1, 2 and 3

$-(b_{13}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{38}^{''})^{(7,7,)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3

$-(b_{36}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)$ are eighth detrition coefficients for category 1, 2 and 3

$-(b_{44}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3



$$\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \left[\begin{array}{ccc} (a'_{44})^{(9)}\boxed{+(a''_{44})^{(9)}(T_{45},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{24})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{36})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)} \end{array}\right] G_{13}$$

$$\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \left[\begin{array}{ccc} (a'_{45})^{(9)}\boxed{+(a''_{45})^{(9)}(T_{45},t)} & \boxed{+(a''_{17})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{21})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{25})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{29})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{33})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{14})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{37})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{41})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)} \end{array}\right] G_{14}$$

$$\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \left[\begin{array}{ccc} (a'_{46})^{(9)}\boxed{+(a''_{46})^{(9)}(T_{37},t)} & \boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{42})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)} \end{array}\right] G_{15}$$

Where $\boxed{+(a''_{44})^{(9)}(T_{45},t)}$, $\boxed{+(a''_{45})^{(9)}(T_{45},t)}$, $\boxed{+(a''_{46})^{(9)}(T_{37},t)}$ are first augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{16})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{17})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{20})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{21})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{24})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{25})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{28})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{29})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{32})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{33})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{13})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a''_{14})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)}$ are Seventh augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)}$ $\boxed{+(a''_{37})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)}$ $\boxed{+(a''_{36})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)}$ are eighth augmentation coefficient for 1,2,3

$\boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a''_{42})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a''_{41})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)}$ are ninth augmentation coefficient for 1,2,3

$\frac{dT_{44}}{dt} =$

$$(b_{44})^{(9)}T_{45} - \left[\begin{array}{ccc} (b_{44}^{'})^{(9)}\boxed{-(b_{44}^{''})^{(9)}(G_{47},t)} & \boxed{-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{array}\right] T_{13}$$

$$\frac{dT_{45}}{dt} = (b_{45})^{(9)}T_{44} - \left[\begin{array}{ccc} (b_{45}^{'})^{(9)}\boxed{-(b_{45}^{''})^{(9)}(G_{47},t)} & \boxed{-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{array}\right] T_{14}$$

$$\frac{dT_{46}}{dt} = (b_{46})^{(9)}T_{45} - \left[\begin{array}{ccc} (b_{46}^{'})^{(9)}\boxed{-(b_{46}^{''})^{(9)}(G_{47},t)} & \boxed{-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{array}\right] T_{15}$$

Where $\boxed{-(b_{44}^{''})^{(9)}(G_{47},t)}$, $\boxed{-(b_{45}^{''})^{(9)}(G_{47},t)}$, $\boxed{-(b_{46}^{''})^{(9)}(G_{47},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)}$ are second detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)}$ are third detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)}$ are fourth detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)}$ are fifth detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)}$ are sixth detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)}$ are seventh detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)}$ are eighth detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$ are ninth detrition coefficients for category 1, 2 and 3

Where we suppose

$(a_i)^{(1)}, (a_i^{'})^{(1)}, (a_i^{''})^{(1)}, (b_i)^{(1)}, (b_i^{'})^{(1)}, (b_i^{''})^{(1)} > 0, \quad i,j = 13,14,15$ *97*

The functions $(a_i^{''})^{(1)}, (b_i^{''})^{(1)}$ are positive continuous increasing and bounded.

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| **Definition of** $(p_i)^{(1)}$, $(r_i)^{(1)}$:<br>$(a_i'')^{(1)}(T_{14},t) \le (p_i)^{(1)} \le (\hat{A}_{13})^{(1)}$<br>$(b_i'')^{(1)}(G,t) \le (r_i)^{(1)} \le (b_i')^{(1)} \le (\hat{B}_{13})^{(1)}$ | |
| $lim_{T_2\to\infty}(a_i'')^{(1)}(T_{14},t) = (p_i)^{(1)}$<br>$\lim_{G\to\infty}(b_i'')^{(1)}(G,t) = (r_i)^{(1)}$<br>**Definition of** $(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}$ :<br>Where $\boxed{(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}}$ are positive constants and $\boxed{i = 13,14,15}$ | 98 |
| They satisfy Lipschitz condition:<br>$\lvert(a_i'')^{(1)}(T_{14}',t) - (a_i'')^{(1)}(T_{14},t)\rvert \le (\hat{k}_{13})^{(1)}\lvert T_{14} - T_{14}'\rvert e^{-(\hat{M}_{13})^{(1)}t}$<br>$\lvert(b_i'')^{(1)}(G',t) - (b_i'')^{(1)}(G,t)\rvert < (\hat{k}_{13})^{(1)}\lVert G - G'\rVert e^{-(\hat{M}_{13})^{(1)}t}$ | 99 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(1)}(T_{14}',t)$ and $(a_i'')^{(1)}(T_{14},t)$ . $(T_{14}',t)$ and $(T_{14},t)$ are points belonging to the interval $\left[(\hat{k}_{13})^{(1)}, (\hat{M}_{13})^{(1)}\right]$. It is to be noted that $(a_i'')^{(1)}(T_{14},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{13})^{(1)} = 1$ then the function $(a_i'')^{(1)}(T_{14},t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}$ :<br>$(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}$, are positive constants<br>$\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}} < 1$ | 100 |
| **Definition of** $(\hat{P}_{13})^{(1)}, (\hat{Q}_{13})^{(1)}$ :<br>There exists two constants $(\hat{P}_{13})^{(1)}$ and $(\hat{Q}_{13})^{(1)}$ which together With $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}, (\hat{A}_{13})^{(1)}$ and $(\hat{B}_{13})^{(1)}$ and the constants $(a_i)^{(1)}, (a_i')^{(1)}, (b_i)^{(1)}, (b_i')^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}, i = 13,14,15,$ satisfy the inequalities<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(a_i)^{(1)} + (a_i')^{(1)} + (\hat{A}_{13})^{(1)} + (\hat{P}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(b_i)^{(1)} + (b_i')^{(1)} + (\hat{B}_{13})^{(1)} + (\hat{Q}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$ | 101 |
| Where we suppose | |
| $(a_i)^{(2)}, (a_i')^{(2)}, (a_i'')^{(2)}, (b_i)^{(2)}, (b_i')^{(2)}, (b_i'')^{(2)} > 0, \quad i,j = 16,17,18$ | |

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|---|---|
| The functions $(a_i'')^{(2)}, (b_i'')^{(2)}$ are positive continuous increasing and bounded. | |
| **Definition of** $(p_i)^{(2)}$, $(r_i)^{(2)}$: | |
| $(a_i'')^{(2)}(T_{17}, t) \leq (p_i)^{(2)} \leq (\hat{A}_{16})^{(2)}$ | 102 |
| $(b_i'')^{(2)}(G_{19}, t) \leq (r_i)^{(2)} \leq (b_i')^{(2)} \leq (\hat{B}_{16})^{(2)}$ | 103 |
| $\lim_{T_2 \to \infty} (a_i'')^{(2)}(T_{17}, t) = (p_i)^{(2)}$ | 104 |
| $\lim_{G \to \infty} (b_i'')^{(2)}((G_{19}), t) = (r_i)^{(2)}$ | 105 |
| **Definition of** $(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}$ :<br>Where $(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}, (p_i)^{(2)}, (r_i)^{(2)}$ are positive constants and $i = 16,17,18$ | 106 |
| They satisfy Lipschitz condition: | |
| $\lvert(a_i'')^{(2)}(T_{17}', t) - (a_i'')^{(2)}(T_{17}, t)\rvert \leq (\hat{k}_{16})^{(2)} \lvert T_{17} - T_{17}'\rvert e^{-(\hat{M}_{16})^{(2)} t}$ | 107 |
| $\lvert(b_i'')^{(2)}((G_{19})', t) - (b_i'')^{(2)}((G_{19}), t)\rvert < (\hat{k}_{16})^{(2)} \lvert\lvert(G_{19}) - (G_{19})'\rvert\rvert e^{-(\hat{M}_{16})^{(2)} t}$ | 108 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(2)}(T_{17}', t)$ and $(a_i'')^{(2)}(T_{17}, t)$ . $(T_{17}', t)$ and $(T_{17}, t)$ are points belonging to the interval $[(\hat{k}_{16})^{(2)}, (\hat{M}_{16})^{(2)}]$ . It is to be noted that $(a_i'')^{(2)}(T_{17}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{16})^{(2)} = 1$ then the function $(a_i'')^{(2)}(T_{17}, t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$ : | |
| $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$, are positive constants<br>$\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}, \frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ | 109 |
| **Definition of** $(\hat{P}_{13})^{(2)}, (\hat{Q}_{13})^{(2)}$ :<br>There exists two constants $(\hat{P}_{16})^{(2)}$ and $(\hat{Q}_{16})^{(2)}$ which together with $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}, (\hat{A}_{16})^{(2)} and (\hat{B}_{16})^{(2)}$ and the constants $(a_i)^{(2)}, (a_i')^{(2)}, (b_i)^{(2)}, (b_i')^{(2)}, (p_i)^{(2)}, (r_i)^{(2)}, i = 16,17,18,$<br>satisfy the inequalities | |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(a_i)^{(2)} + (a_i')^{(2)} + (\hat{A}_{16})^{(2)} + (\hat{P}_{16})^{(2)} (\hat{k}_{16})^{(2)}] < 1$ | 110 |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(b_i)^{(2)} + (b_i')^{(2)} + (\hat{B}_{16})^{(2)} + (\hat{Q}_{16})^{(2)} (\hat{k}_{16})^{(2)}] < 1$ | 111 |

| | |
|---|---|
| Where we suppose | |
| $(a_i)^{(3)}, (a_i^{'})^{(3)}, (a_i^{''})^{(3)}, (b_i)^{(3)}, (b_i^{'})^{(3)}, (b_i^{''})^{(3)} > 0, \quad i,j = 20,21,22$<br><br>The functions $(a_i^{''})^{(3)}, (b_i^{''})^{(3)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(3)}$, $(\mathrm{r_i})^{(3)}$:<br><br>$(a_i^{''})^{(3)}(T_{21},t) \le (p_i)^{(3)} \le (\hat{A}_{20})^{(3)}$<br><br>$(b_i^{''})^{(3)}(G_{23},t) \le (r_i)^{(3)} \le (b_i^{'})^{(3)} \le (\hat{B}_{20})^{(3)}$ | 112 |
| $lim_{T_2\to\infty}(a_i^{''})^{(3)}(T_{21},t) = (p_i)^{(3)}$<br><br>$\lim_{G\to\infty}(b_i^{''})^{(3)}(G_{23},t) = (r_i)^{(3)}$<br><br>**Definition of** $(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}$ :<br><br>Where $\boxed{(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}}$ are positive constants and $\boxed{i = 20,21,22}$ | 113 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(3)}(T_{21}^{'},t) - (a_i^{''})^{(3)}(T_{21},t)\rvert \le (\hat{k}_{20})^{(3)}\lvert T_{21} - T_{21}^{'}\rvert e^{-(\hat{M}_{20})^{(3)}t}$<br><br>$\lvert(b_i^{''})^{(3)}(G_{23}^{'},t) - (b_i^{''})^{(3)}(G_{23},t)\rvert < (\hat{k}_{20})^{(3)}\lVert G_{23} - G_{23}^{'}\rVert e^{-(\hat{M}_{20})^{(3)}t}$ | 114 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(3)}(T_{21}^{'},t)$ and$(a_i^{''})^{(3)}(T_{21},t)$ . $(T_{21}^{'},t)$ And $(T_{21},t)$ are points belonging to the interval $\left[(\hat{k}_{20})^{(3)}, (\hat{M}_{20})^{(3)}\right]$. It is to be noted that $(a_i^{''})^{(3)}(T_{21},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{20})^{(3)} = 1$ then the function $(a_i^{''})^{(3)}(T_{21},t)$, the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$ :<br><br>$(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$, are positive constants<br>$\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}, \frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}} < 1$ | 115 |
| There exists two constants There exists two constants $(\hat{P}_{20})^{(3)}$ and $(\hat{Q}_{20})^{(3)}$ which together with $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}, (\hat{A}_{20})^{(3)} and\ (\hat{B}_{20})^{(3)}$ and the constants $(a_i)^{(3)}, (a_i^{'})^{(3)}, (b_i)^{(3)}, (b_i^{'})^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}, i = 20,21,22,$<br>satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(a_i)^{(3)} + (a_i^{'})^{(3)} + (\hat{A}_{20})^{(3)} + (\hat{P}_{20})^{(3)}(\hat{k}_{20})^{(3)}] < 1$<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(b_i)^{(3)} + (b_i^{'})^{(3)} + (\hat{B}_{20})^{(3)} + (\hat{Q}_{20})^{(3)}(\hat{k}_{20})^{(3)}] < 1$ | 116 |

| | |
|---|---|
| Where we suppose | |
| $(a_i)^{(4)}, (a_i^{'})^{(4)}, (a_i^{''})^{(4)}, (b_i)^{(4)}, (b_i^{'})^{(4)}, (b_i^{''})^{(4)} > 0, \quad i,j = 24{,}25{,}26$<br><br>The functions $(a_i^{''})^{(4)}, (b_i^{''})^{(4)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(4)}$, $(r_i)^{(4)}$:<br><br>$(a_i^{''})^{(4)}(T_{25}, t) \leq (p_i)^{(4)} \leq (\hat{A}_{24})^{(4)}$<br><br>$(b_i^{''})^{(4)}\big((G_{27}), t\big) \leq (r_i)^{(4)} \leq (b_i^{'})^{(4)} \leq (\hat{B}_{24})^{(4)}$ | 117 |
| $lim_{T_2 \to \infty} (a_i^{''})^{(4)}(T_{25}, t) = (p_i)^{(4)}$<br><br>$\lim_{G \to \infty} (b_i^{''})^{(4)}\big((G_{27}), t\big) = (r_i)^{(4)}$<br><br>**Definition of** $(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}$ :<br><br>Where $\boxed{(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}, (p_i)^{(4)}, (r_i)^{(4)}}$ are positive constants and $\boxed{i = 24{,}25{,}26}$ | 118 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(4)}(T_{25}^{'}, t) - (a_i^{''})^{(4)}(T_{25}, t)\rvert \leq (\hat{k}_{24})^{(4)} \lvert T_{25} - T_{25}^{'}\rvert e^{-(\hat{M}_{24})^{(4)} t}$<br><br>$\lvert(b_i^{''})^{(4)}((G_{27})^{'}, t) - (b_i^{''})^{(4)}\big((G_{27}), t\big)\rvert < (\hat{k}_{24})^{(4)} \lvert\lvert(G_{27}) - (G_{27})^{'}\rvert\rvert e^{-(\hat{M}_{24})^{(4)} t}$ | 119 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(4)}(T_{25}^{'}, t)$ and $(a_i^{''})^{(4)}(T_{25}, t)$ . $(T_{25}^{'}, t)$ and $(T_{25}, t)$ are points belonging to the interval $\big[(\hat{k}_{24})^{(4)}, (\hat{M}_{24})^{(4)}\big]$. It is to be noted that $(a_i^{''})^{(4)}(T_{25}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{24})^{(4)} = 1$ then the function $(a_i^{''})^{(4)}(T_{25}, t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$ :<br><br>$(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$, are positive constants<br><br>$\frac{(a_i)^{(4)}}{(\hat{M}_{24})^{(4)}}, \frac{(b_i)^{(4)}}{(\hat{M}_{24})^{(4)}} < 1$ | 120 |
| **Definition of** $(\hat{P}_{24})^{(4)}, (\hat{Q}_{24})^{(4)}$ :<br><br>There exists two constants $(\hat{P}_{24})^{(4)}$ and $(\hat{Q}_{24})^{(4)}$ which together with $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}, (\hat{A}_{24})^{(4)} and\ (\hat{B}_{24})^{(4)}$ and the constants $(a_i)^{(4)}, (a_i^{'})^{(4)}, (b_i)^{(4)}, (b_i^{'})^{(4)}, (p_i)^{(4)}, (r_i)^{(4)}, i = 24{,}25{,}26$, satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{24})^{(4)}}[(a_i)^{(4)} + (a_i^{'})^{(4)} + (\hat{A}_{24})^{(4)} + (\hat{P}_{24})^{(4)} (\hat{k}_{24})^{(4)}] < 1$ | 121 |

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| $\frac{1}{(\hat{M}_{24})^{(4)}}[\ (b_i)^{(4)} + (b_i^{'})^{(4)} +\ (\hat{B}_{24})^{(4)} + (\hat{Q}_{24})^{(4)}\ (\hat{k}_{24})^{(4)}] < 1$ | |
| Where we suppose | |
| $(a_i)^{(5)}, (a_i^{'})^{(5)}, (a_i^{''})^{(5)}, (b_i)^{(5)}, (b_i^{'})^{(5)}, (b_i^{''})^{(5)} > 0, \quad i,j = 28,29,30$<br><br>The functions $(a_i^{''})^{(5)}, (b_i^{''})^{(5)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(5)}$, $(r_i)^{(5)}$:<br><br>$(a_i^{''})^{(5)}(T_{29}, t) \leq (p_i)^{(5)} \leq (\hat{A}_{28})^{(5)}$<br><br>$(b_i^{''})^{(5)}\big((G_{31}), t\big) \leq\ (r_i)^{(5)} \leq (b_i^{'})^{(5)} \leq (\hat{B}_{28})^{(5)}$ | 122 |
| $lim_{T_2 \to \infty}(a_i^{''})^{(5)}\,(T_{29}, t) = (p_i)^{(5)}$<br><br>$\lim_{G \to \infty}(b_i^{''})^{(5)}\,(G_{31}, t) =\ (r_i)^{(5)}$<br><br>**Definition of** $(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}$ :<br><br>Where $\boxed{(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}, (p_i)^{(5)},\ (r_i)^{(5)}}$ are positive constants and $\boxed{i = 28,29,30}$ | 123 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(5)}(T_{29}^{'}, t) - (a_i^{''})^{(5)}(T_{29}, t)\rvert \leq (\hat{k}_{28})^{(5)}\lvert T_{29} - T_{29}^{'}\rvert e^{-(\hat{M}_{28})^{(5)}t}$<br><br>$\lvert(b_i^{''})^{(5)}((G_{31})^{'}, t) - (b_i^{''})^{(5)}\big((G_{31}), t\big)\rvert < (\hat{k}_{28})^{(5)}\lvert\lvert(G_{31}) - (G_{31})^{'}\rvert\rvert e^{-(\hat{M}_{28})^{(5)}t}$ | 124 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(5)}(T_{29}^{'}, t)$ and $(a_i^{''})^{(5)}(T_{29}, t)$ . $(T_{29}^{'}, t)$ and $(T_{29}, t)$ are points belonging to the interval $\big[(\hat{k}_{28})^{(5)}, (\hat{M}_{28})^{(5)}\big]$. It is to be noted that $(a_i^{''})^{(5)}(T_{29}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{28})^{(5)} = 1$ then the function $(a_i^{''})^{(5)}(T_{29}, t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}$ :<br><br>$(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}$, are positive constants<br><br>$\frac{(a_i)^{(5)}}{(\hat{M}_{28})^{(5)}}\ , \frac{(b_i)^{(5)}}{(\hat{M}_{28})^{(5)}} < 1$ | 125 |
| **Definition of** $(\hat{P}_{28})^{(5)}, (\hat{Q}_{28})^{(5)}$ :<br><br>There exists two constants $(\hat{P}_{28})^{(5)}$ and $(\hat{Q}_{28})^{(5)}$ which together with $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}, (\hat{A}_{28})^{(5)} and\ (\hat{B}_{28})^{(5)}$ and the constants $(a_i)^{(5)}, (a_i^{'})^{(5)}, (b_i)^{(5)}, (b_i^{'})^{(5)}, (p_i)^{(5)},\ (r_i)^{(5)}, i = 28,29,30,$ satisfy the inequalities | 126 |

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| $\frac{1}{(\hat{M}_{28})^{(5)}}[(a_i)^{(5)} + (a_i^{'})^{(5)} + (\hat{A}_{28})^{(5)} + (\hat{P}_{28})^{(5)}(\hat{k}_{28})^{(5)}] < 1$<br><br>$\frac{1}{(\hat{M}_{28})^{(5)}}[(b_i)^{(5)} + (b_i^{'})^{(5)} + (\hat{B}_{28})^{(5)} + (\hat{Q}_{28})^{(5)}(\hat{k}_{28})^{(5)}] < 1$ | |
| Where we suppose | |
| $(a_i)^{(6)}, (a_i^{'})^{(6)}, (a_i^{''})^{(6)}, (b_i)^{(6)}, (b_i^{'})^{(6)}, (b_i^{''})^{(6)} > 0, \quad i,j = 32,33,34$<br><br>The functions $(a_i^{''})^{(6)}, (b_i^{''})^{(6)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(6)}$, $(r_i)^{(6)}$:<br><br>$(a_i^{''})^{(6)}(T_{33}, t) \leq (p_i)^{(6)} \leq (\hat{A}_{32})^{(6)}$<br><br>$(b_i^{''})^{(6)}((G_{35}), t) \leq (r_i)^{(6)} \leq (b_i^{'})^{(6)} \leq (\hat{B}_{32})^{(6)}$ | 127 |
| $lim_{T_2 \to \infty}(a_i^{''})^{(6)}(T_{33}, t) = (p_i)^{(6)}$<br><br>$\lim_{G \to \infty}(b_i^{''})^{(6)}((G_{35}), t) = (r_i)^{(6)}$<br><br>**Definition of** $(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}$ :<br><br>Where $\boxed{(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}, (p_i)^{(6)}, (r_i)^{(6)}}$ are positive constants and $\boxed{i = 32,33,34}$ | 128 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(6)}(T_{33}^{'}, t) - (a_i^{''})^{(6)}(T_{33}, t)\rvert \leq (\hat{k}_{32})^{(6)}\lvert T_{33} - T_{33}^{'}\rvert e^{-(\hat{M}_{32})^{(6)}t}$<br><br>$\lvert(b_i^{''})^{(6)}((G_{35})^{'}, t) - (b_i^{''})^{(6)}((G_{35}), t)\rvert < (\hat{k}_{32})^{(6)}\lvert\lvert(G_{35}) - (G_{35})^{'}\rvert\rvert e^{-(\hat{M}_{32})^{(6)}t}$ | |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(6)}(T_{33}^{'}, t)$ and$(a_i^{''})^{(6)}(T_{33}, t)$ . $(T_{33}^{'}, t)$ and $(T_{33}, t)$ are points belonging to the interval $[(\hat{k}_{32})^{(6)}, (\hat{M}_{32})^{(6)}]$. It is to be noted that $(a_i^{''})^{(6)}(T_{33}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{32})^{(6)} = 1$ then the function $(a_i^{''})^{(6)}(T_{33}, t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}$ :<br><br>$(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}$, are positive constants<br><br>$\frac{(a_i)^{(6)}}{(\hat{M}_{32})^{(6)}}, \frac{(b_i)^{(6)}}{(\hat{M}_{32})^{(6)}} < 1$ | 129 |
| **Definition of** $(\hat{P}_{32})^{(6)}, (\hat{Q}_{32})^{(6)}$ :<br><br>There exists two constants $(\hat{P}_{32})^{(6)}$ and $(\hat{Q}_{32})^{(6)}$ which together with $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}, (\hat{A}_{32})^{(6)} and\ (\hat{B}_{32})^{(6)}$ and the constants | 130 |

$(a_i)^{(6)}, (a_i^{'})^{(6)}, (b_i)^{(6)}, (b_i^{'})^{(6)}, (p_i)^{(6)}, \ (r_i)^{(6)}, i = 32,33,34,$
satisfy the inequalities

$$\frac{1}{(\hat{M}_{32})^{(6)}}[(a_i)^{(6)} + (a_i^{'})^{(6)} + (\hat{A}_{32})^{(6)} + (\hat{P}_{32})^{(6)} (\hat{k}_{32})^{(6)}] < 1$$

$$\frac{1}{(\hat{M}_{32})^{(6)}}[(b_i)^{(6)} + (b_i^{'})^{(6)} + (\hat{B}_{32})^{(6)} + (\hat{Q}_{32})^{(6)} (\hat{k}_{32})^{(6)}] < 1$$

Where we suppose

(G) $(a_i)^{(7)}, (a_i^{'})^{(7)}, (a_i^{''})^{(7)}, (b_i)^{(7)}, (b_i^{'})^{(7)}, (b_i^{''})^{(7)} > 0, \quad i, j = 36,37,38$ 131

(H) The functions $(a_i^{''})^{(7)}, (b_i^{''})^{(7)}$ are positive continuous increasing and bounded.

**Definition of** $(p_i)^{(7)}, \ (r_i)^{(7)}$:

$$(a_i^{''})^{(7)}(T_{37}, t) \leq (p_i)^{(7)} \leq (\hat{A}_{36})^{(7)}$$

$$(b_i^{''})^{(7)}(G_{39}, t) \leq (r_i)^{(7)} \leq (b_i^{'})^{(7)} \leq (\hat{B}_{36})^{(7)}$$

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*(I)* $\lim_{T_2 \to \infty} (a_i^{''})^{(7)}(T_{37}, t) = (p_i)^{(7)}$

*(J)*

$\lim_{G \to \infty} (b_i^{''})^{(7)}((G_{39}), t) = (r_i)^{(7)}$

**Definition of** $(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}$ :

Where $\boxed{(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}, (p_i)^{(7)}, \ (r_i)^{(7)}}$ are positive constants and $\boxed{i = 36,37,38}$

They satisfy Lipschitz condition: 133

$$|(a_i^{''})^{(7)}(T_{37}^{'}, t) - (a_i^{''})^{(7)}(T_{37}, t)| \leq (\hat{k}_{36})^{(7)} |T_{37} - T_{37}^{'}| e^{-(\hat{M}_{36})^{(7)} t}$$

$$|(b_i^{''})^{(7)}((G_{39})^{'}, t) - (b_i^{''})^{(7)}((G_{39}), t)| < (\hat{k}_{36})^{(7)} ||(G_{39}) - (G_{39})^{'}|| e^{-(\hat{M}_{36})^{(7)} t}$$

With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(7)}(T_{37}^{'}, t)$ and $(a_i^{''})^{(7)}(T_{37}, t)$ . $(T_{37}^{'}, t)$ and $(T_{37}, t)$ are points belonging to the interval $[(\hat{k}_{36})^{(7)}, (\hat{M}_{36})^{(7)}]$ . It is to be noted that $(a_i^{''})^{(7)}(T_{37}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{36})^{(7)} = 1$ then the function $(a_i^{''})^{(7)}(T_{37}, t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

**Definition of** $(\hat{M}_{36})^{(7)}, (\hat{k}_{36})^{(7)}$ : 134

(K) $(\hat{M}_{36})^{(7)}, (\hat{k}_{36})^{(7)},$ are positive constants

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| $\frac{(a_i)^{(7)}}{(\hat{M}_{36})^{(7)}}, \frac{(b_i)^{(7)}}{(\hat{M}_{36})^{(7)}} < 1$ | |
| **Definition of** $(\hat{P}_{36})^{(7)}, (\hat{Q}_{36})^{(7)}$ :<br><br>(L) There exists two constants $(\hat{P}_{36})^{(7)}$ and $(\hat{Q}_{36})^{(7)}$ which together with $(\hat{M}_{36})^{(7)}, (\hat{k}_{36})^{(7)}, (\hat{A}_{36})^{(7)} and\ (\hat{B}_{36})^{(7)}$ and the constants $(a_i)^{(7)}, (a_i^{'})^{(7)}, (b_i)^{(7)}, (b_i^{'})^{(7)}, (p_i)^{(7)}, (r_i)^{(7)}, i = 36,37,38$, satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{36})^{(7)}}[(a_i)^{(7)} + (a_i^{'})^{(7)} + (\hat{A}_{36})^{(7)} + (\hat{P}_{36})^{(7)} (\hat{k}_{36})^{(7)}] < 1$<br><br>$\frac{1}{(\hat{M}_{36})^{(7)}}[(b_i)^{(7)} + (b_i^{'})^{(7)} + (\hat{B}_{36})^{(7)} + (\hat{Q}_{36})^{(7)} (\hat{k}_{36})^{(7)}] < 1$ | 135 |
| Where we suppose | |
| $(a_i)^{(8)}, (a_i^{'})^{(8)}, (a_i^{''})^{(8)}, (b_i)^{(8)}, (b_i^{'})^{(8)}, (b_i^{''})^{(8)} > 0, \quad i,j = 40,41,42$ | 136 |
| The functions $(a_i^{''})^{(8)}, (b_i^{''})^{(8)}$ are positive continuous increasing and bounded | |
| **Definition of** $(p_i)^{(8)}, (r_i)^{(8)}$: | 137 |
| $(a_i^{''})^{(8)}(T_{41}, t) \leq (p_i)^{(8)} \leq (\hat{A}_{40})^{(8)}$ | 138 |
| $(b_i^{''})^{(8)}((G_{43}), t) \leq (r_i)^{(8)} \leq (b_i^{'})^{(8)} \leq (\hat{B}_{40})^{(8)}$ | 139 |
| $\lim_{T_2 \to \infty} (a_i^{''})^{(8)}(T_{41}, t) = (p_i)^{(8)}$ | 140 |
| $\lim_{G \to \infty} (b_i^{''})^{(8)}((G_{43}), t) = (r_i)^{(8)}$ | 141 |
| **Definition of** $(\hat{A}_{40})^{(8)}, (\hat{B}_{40})^{(8)}$ :<br><br>Where $\boxed{(\hat{A}_{40})^{(8)}, (\hat{B}_{40})^{(8)}, (p_i)^{(8)}, (r_i)^{(8)}}$ are positive constants and $\boxed{i = 40,41,42}$ | |
| They satisfy Lipschitz condition: | |
| $\lvert(a_i^{''})^{(8)}(T_{41}^{'}, t) - (a_i^{''})^{(8)}(T_{41}, t)\rvert \leq (\hat{k}_{40})^{(8)} \lvert T_{41} - T_{41}^{'}\rvert e^{-(\hat{M}_{40})^{(8)} t}$ | 142 |
| $\lvert(b_i^{''})^{(8)}((G_{43})', t) - (b_i^{''})^{(8)}((G_{43}), t)\rvert < (\hat{k}_{40})^{(8)} \lvert\lvert(G_{43}) - (G_{43})'\rvert\rvert e^{-(\hat{M}_{40})^{(8)} t}$ | 143 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(8)}(T_{41}^{'}, t)$ and $(a_i^{''})^{(8)}(T_{41}, t)$ . $(T_{41}^{'}, t)$ and $(T_{41}, t)$ are points belonging to the interval $[(\hat{k}_{40})^{(8)}, (\hat{M}_{40})^{(8)}]$. It is to be noted that $(a_i^{''})^{(8)}(T_{41}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{40})^{(8)} = 1$ then the function $(a_i^{''})^{(8)}(T_{41}, t)$ , the **first augmentation coefficient** attributable to the system, would | |

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| be absolutely continuous. | |
| **Definition of** $(\widehat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}$ : | |
| $(\widehat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)},$ are positive constants | |
| $\frac{(a_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}$ , $\frac{(b_i)^{(8)}}{(\widehat{M}_{40})^{(8)}} < 1$ | 144 |
| **Definition of** $(\hat{P}_{40})^{(8)}, (\hat{Q}_{40})^{(8)}$ :<br>There exists two constants $(\hat{P}_{40})^{(8)}$ and $(\hat{Q}_{40})^{(8)}$ which together with $(\widehat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}, (\hat{A}_{40})^{(8)}$ $(\hat{B}_{40})^{(8)}$ and the constants $(a_i)^{(8)}, (a_i')^{(8)}, (b_i)^{(8)}, (b_i')^{(8)}, (p_i)^{(8)}, (r_i)^{(8)}, i = 40,41,42,$<br>Satisfy the inequalities | |
| $\frac{1}{(\widehat{M}_{40})^{(8)}}[(a_i)^{(8)} + (a_i')^{(8)} + (\hat{A}_{40})^{(8)} + (\hat{P}_{40})^{(8)} (\hat{k}_{40})^{(8)}] < 1$ | 145 |
| $\frac{1}{(\widehat{M}_{40})^{(8)}}[(b_i)^{(8)} + (b_i')^{(8)} + (\hat{B}_{40})^{(8)} + (\hat{Q}_{40})^{(8)} (\hat{k}_{40})^{(8)}] < 1$ | 146 |
| Where we suppose | |
| $(a_i)^{(9)}, (a_i')^{(9)}, (a_i'')^{(9)}, (b_i)^{(9)}, (b_i')^{(9)}, (b_i'')^{(9)} > 0, \quad i, j = 44,45,46$<br>The functions $(a_i'')^{(9)}, (b_i'')^{(9)}$ are positive continuous increasing and bounded.<br>**Definition of** $(p_i)^{(9)}$, $(r_i)^{(9)}$:<br>$(a_i'')^{(9)}(T_{45}, t) \le (p_i)^{(9)} \le (\hat{A}_{44})^{(9)}$<br>$(b_i'')^{(9)}(G_{47}, t) \le (r_i)^{(9)} \le (b_i')^{(9)} \le (\hat{B}_{44})^{(9)}$ | 146 A |
| $\lim_{T_2 \to \infty} (a_i'')^{(9)}(T_{45}, t) = (p_i)^{(9)}$<br>$\lim_{G \to \infty} (b_i'')^{(9)}(G_{47}, t) = (r_i)^{(9)}$<br>**Definition of** $(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}$ :<br>Where $\boxed{(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}, (p_i)^{(9)}, (r_i)^{(9)}}$ are positive constants and $\boxed{i = 44,45,46}$ | |
| They satisfy Lipschitz condition:<br>$\lvert(a_i'')^{(9)}(T_{45}', t) - (a_i'')^{(9)}(T_{45}, t)\rvert \le (\hat{k}_{44})^{(9)} \lvert T_{45} - T_{45}' \rvert e^{-(\widehat{M}_{44})^{(9)} t}$<br>$\lvert(b_i'')^{(9)}((G_{47})', t) - (b_i'')^{(9)}((G_{47}), t)\rvert < (\hat{k}_{44})^{(9)} \lvert\lvert(G_{47}) - (G_{47})'\rvert\rvert e^{-(\widehat{M}_{44})^{(9)} t}$ | |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(9)}(T_{45}', t)$ | |

and $(a_i'')^{(9)}(T_{45},t)$ . $(T_{45}',t)$ and $(T_{45},t)$ are points belonging to the interval $\left[(\hat{k}_{44})^{(9)},(\widehat{M}_{44})^{(9)}\right]$. It is to be noted that $(a_i'')^{(9)}(T_{45},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\widehat{M}_{44})^{(9)}=1$ then the function $(a_i'')^{(9)}(T_{45},t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

**Definition of** $(\widehat{M}_{44})^{(9)},(\hat{k}_{44})^{(9)}$ :

$(\widehat{M}_{44})^{(9)},(\hat{k}_{44})^{(9)},$ are positive constants

$$\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\ ,\frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}<1$$

**Definition of** $(\hat{P}_{44})^{(9)},(\hat{Q}_{44})^{(9)}$ :

There exists two constants $(\hat{P}_{44})^{(9)}$ and $(\hat{Q}_{44})^{(9)}$ which together with $(\widehat{M}_{44})^{(9)},(\hat{k}_{44})^{(9)},(\hat{A}_{44})^{(9)} and\ (\hat{B}_{44})^{(9)}$ and the constants $(a_i)^{(9)},(a_i')^{(9)},(b_i)^{(9)},(b_i')^{(9)},(p_i)^{(9)},\ (r_i)^{(9)},i=44,45,46,$ satisfy the inequalities

$$\frac{1}{(\widehat{M}_{44})^{(9)}}[(a_i)^{(9)}+(a_i')^{(9)}+\ (\hat{A}_{44})^{(9)}+\ (\hat{P}_{44})^{(9)}\ (\hat{k}_{44})^{(9)}]<1$$

$$\frac{1}{(\widehat{M}_{44})^{(9)}}[\ (b_i)^{(9)}+(b_i')^{(9)}+\ (\hat{B}_{44})^{(9)}+\ (\hat{Q}_{44})^{(9)}\ \ (\hat{k}_{44})^{(9)}]<1$$

**Theorem 1:** if the conditions above are fulfilled, there exists a solution satisfying the conditions 147

Definition of $G_i(0)\ ,T_i(0)$ :

$$G_i(t)\le\ (\hat{P}_{13})^{(1)}e^{(\widehat{M}_{13})^{(1)}t}\ ,\ \boxed{G_i(0)=G_i^0>0}$$

$$T_i(t)\le\ (\hat{Q}_{13})^{(1)}e^{(\widehat{M}_{13})^{(1)}t}\ ,\ \boxed{T_i(0)=T_i^0>0}$$

**Theorem 2 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions 148

**Definition of** $G_i(0)\ ,T_i(0)$

$$G_i(t)\le\ (\hat{P}_{16})^{(2)}e^{(\widehat{M}_{16})^{(2)}t}\ ,\quad G_i(0)=G_i^0>0$$

$$T_i(t)\le\ (\hat{Q}_{16})^{(2)}e^{(\widehat{M}_{16})^{(2)}t}\ ,\quad T_i(0)=T_i^0>0$$

**Theorem 3 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions 149

$$G_i(t)\le\ (\hat{P}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}\ ,\quad G_i(0)=G_i^0>0$$

$$T_i(t)\le\ (\hat{Q}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}\ ,\quad T_i(0)=T_i^0>0$$

**Theorem 4 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions 150

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| **Definition of** $G_i(0)\,, T_i(0)$ :<br>$G_i(t) \le \left(\hat{P}_{24}\right)^{(4)} e^{(\hat{M}_{24})^{(4)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{24}\right)^{(4)} e^{(\hat{M}_{24})^{(4)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | |
| **Theorem 5 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0)\,, T_i(0)$ :<br>$G_i(t) \le \left(\hat{P}_{28}\right)^{(5)} e^{(\hat{M}_{28})^{(5)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{28}\right)^{(5)} e^{(\hat{M}_{28})^{(5)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 151 |
| **Theorem 6 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0)\,, T_i(0)$ :<br>$G_i(t) \le \left(\hat{P}_{32}\right)^{(6)} e^{(\hat{M}_{32})^{(6)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{32}\right)^{(6)} e^{(\hat{M}_{32})^{(6)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 152 |
| **Theorem 7:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0)\,, T_i(0)$ :<br>$G_i(t) \le \left(\hat{P}_{36}\right)^{(7)} e^{(\hat{M}_{36})^{(7)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{36}\right)^{(7)} e^{(\hat{M}_{36})^{(7)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 |
| **Theorem 8:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0)\,, T_i(0)$ :<br>$G_i(t) \le \left(\hat{P}_{40}\right)^{(8)} e^{(\hat{M}_{40})^{(8)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{40}\right)^{(8)} e^{(\hat{M}_{40})^{(8)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 A |
| **Theorem 9:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0)\,, T_i(0)$ :<br>$G_i(t) \le \left(\hat{P}_{44}\right)^{(9)} e^{(\hat{M}_{44})^{(9)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$ | 153 B |

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| $T_i(t) \leq (\hat{Q}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | |
| **Proof:** Consider operator $\mathcal{A}^{(1)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 154 |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{13})^{(1)}$ , $T_i^0 \leq (\hat{Q}_{13})^{(1)}$, | 155 |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ | 156 |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ | 157 |
| By<br>$\bar{G}_{13}(t) = G_{13}^0 + \int_0^t \left[ (a_{13})^{(1)} G_{14}(s_{(13)}) - \left( (a_{13}')^{(1)} + a_{13}'')^{(1)} (T_{14}(s_{(13)}), s_{(13)}) \right) G_{13}(s_{(13)}) \right] ds_{(13)}$ | 158 |
| $\bar{G}_{14}(t) = G_{14}^0 + \int_0^t \left[ (a_{14})^{(1)} G_{13}(s_{(13)}) - \left( (a_{14}')^{(1)} + (a_{14}'')^{(1)} (T_{14}(s_{(13)}), s_{(13)}) \right) G_{14}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{G}_{15}(t) = G_{15}^0 + \int_0^t \left[ (a_{15})^{(1)} G_{14}(s_{(13)}) - \left( (a_{15}')^{(1)} + (a_{15}'')^{(1)} (T_{14}(s_{(13)}), s_{(13)}) \right) G_{15}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{T}_{13}(t) = T_{13}^0 + \int_0^t \left[ (b_{13})^{(1)} T_{14}(s_{(13)}) - \left( (b_{13}')^{(1)} - (b_{13}'')^{(1)} (G(s_{(13)}), s_{(13)}) \right) T_{13}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{T}_{14}(t) = T_{14}^0 + \int_0^t \left[ (b_{14})^{(1)} T_{13}(s_{(13)}) - \left( (b_{14}')^{(1)} - (b_{14}'')^{(1)} (G(s_{(13)}), s_{(13)}) \right) T_{14}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{T}_{15}(t) = T_{15}^0 + \int_0^t \left[ (b_{15})^{(1)} T_{14}(s_{(13)}) - \left( (b_{15}')^{(1)} - (b_{15}'')^{(1)} (G(s_{(13)}), s_{(13)}) \right) T_{15}(s_{(13)}) \right] ds_{(13)}$<br>Where $s_{(13)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(2)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 159 |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{16})^{(2)}$ , $T_i^0 \leq (\hat{Q}_{16})^{(2)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ | |
| By<br>$\bar{G}_{16}(t) = G_{16}^0 + \int_0^t \left[ (a_{16})^{(2)} G_{17}(s_{(16)}) - \left( (a_{16}')^{(2)} + a_{16}'')^{(2)} (T_{17}(s_{(16)}), s_{(16)}) \right) G_{16}(s_{(16)}) \right] ds_{(16)}$ | 160 |
| $\bar{G}_{17}(t) = G_{17}^0 + \int_0^t \left[ (a_{17})^{(2)} G_{16}(s_{(16)}) - \left( (a_{17}')^{(2)} + (a_{17}'')^{(2)} (T_{17}(s_{(16)}), s_{(17)}) \right) G_{17}(s_{(16)}) \right] ds_{(16)}$ | |

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| $\bar{G}_{18}(t) = G_{18}^0 + \int_0^t \left[(a_{18})^{(2)} G_{17}(s_{(16)}) - \left((a_{18}')^{(2)} + (a_{18}'')^{(2)}(T_{17}(s_{(16)}), s_{(16)})\right) G_{18}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{T}_{16}(t) = T_{16}^0 + \int_0^t \left[(b_{16})^{(2)} T_{17}(s_{(16)}) - \left((b_{16}')^{(2)} - (b_{16}'')^{(2)}(G_{19}(s_{(16)}), s_{(16)})\right) T_{16}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{T}_{17}(t) = T_{17}^0 + \int_0^t \left[(b_{17})^{(2)} T_{16}(s_{(16)}) - \left((b_{17}')^{(2)} - (b_{17}'')^{(2)}(G_{19}(s_{(16)}), s_{(16)})\right) T_{17}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{T}_{18}(t) = T_{18}^0 + \int_0^t \left[(b_{18})^{(2)} T_{17}(s_{(16)}) - \left((b_{18}')^{(2)} - (b_{18}'')^{(2)}(G_{19}(s_{(16)}), s_{(16)})\right) T_{18}(s_{(16)})\right] ds_{(16)}$<br>Where $s_{(16)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(3)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{20})^{(3)}$, $T_i^0 \le (\hat{Q}_{20})^{(3)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ | |
| By<br>$\bar{G}_{20}(t) = G_{20}^0 + \int_0^t \left[(a_{20})^{(3)} G_{21}(s_{(20)}) - \left((a_{20}')^{(3)} + a_{20}'')^{(3)}(T_{21}(s_{(20)}), s_{(20)})\right) G_{20}(s_{(20)})\right] ds_{(20)}$ | 161 |
| $\bar{G}_{21}(t) = G_{21}^0 + \int_0^t \left[(a_{21})^{(3)} G_{20}(s_{(20)}) - \left((a_{21}')^{(3)} + (a_{21}'')^{(3)}(T_{21}(s_{(20)}), s_{(20)})\right) G_{21}(s_{(20)})\right] ds_{(20)}$ | |
| $\bar{G}_{22}(t) = G_{22}^0 + \int_0^t \left[(a_{22})^{(3)} G_{21}(s_{(20)}) - \left((a_{22}')^{(3)} + (a_{22}'')^{(3)}(T_{21}(s_{(20)}), s_{(20)})\right) G_{22}(s_{(20)})\right] ds_{(20)}$ | |
| $\bar{T}_{20}(t) = T_{20}^0 + \int_0^t \left[(b_{20})^{(3)} T_{21}(s_{(20)}) - \left((b_{20}')^{(3)} - (b_{20}'')^{(3)}(G_{23}(s_{(20)}), s_{(20)})\right) T_{20}(s_{(20)})\right] ds_{(20)}$ | |
| $\bar{T}_{21}(t) = T_{21}^0 + \int_0^t \left[(b_{21})^{(3)} T_{20}(s_{(20)}) - \left((b_{21}')^{(3)} - (b_{21}'')^{(3)}(G_{23}(s_{(20)}), s_{(20)})\right) T_{21}(s_{(20)})\right] ds_{(20)}$ | |
| $\bar{T}_{22}(t) = T_{22}^0 + \int_0^t \left[(b_{22})^{(3)} T_{21}(s_{(20)}) - \left((b_{22}')^{(3)} - (b_{22}'')^{(3)}(G_{23}(s_{(20)}), s_{(20)})\right) T_{22}(s_{(20)})\right] ds_{(20)}$<br>Where $s_{(20)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:** Consider operator $\mathcal{A}^{(4)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{24})^{(4)}$, $T_i^0 \le (\hat{Q}_{24})^{(4)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ | |

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| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ | |
| By<br>$\bar{G}_{24}(t) = G_{24}^0 + \int_0^t \left[ (a_{24})^{(4)} G_{25}(s_{(24)}) - \left( (a'_{24})^{(4)} + a''_{24})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{24}(s_{(24)}) \right] ds_{(24)}$ | 162 |
| $\bar{G}_{25}(t) = G_{25}^0 + \int_0^t \left[ (a_{25})^{(4)} G_{24}(s_{(24)}) - \left( (a'_{25})^{(4)} + (a''_{25})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{25}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{G}_{26}(t) = G_{26}^0 + \int_0^t \left[ (a_{26})^{(4)} G_{25}(s_{(24)}) - \left( (a'_{26})^{(4)} + (a''_{26})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{26}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{24}(t) = T_{24}^0 + \int_0^t \left[ (b_{24})^{(4)} T_{25}(s_{(24)}) - \left( (b'_{24})^{(4)} - (b''_{24})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{24}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{25}(t) = T_{25}^0 + \int_0^t \left[ (b_{25})^{(4)} T_{24}(s_{(24)}) - \left( (b'_{25})^{(4)} - (b''_{25})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{25}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{26}(t) = T_{26}^0 + \int_0^t \left[ (b_{26})^{(4)} T_{25}(s_{(24)}) - \left( (b'_{26})^{(4)} - (b''_{26})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{26}(s_{(24)}) \right] ds_{(24)}$<br>Where $s_{(24)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:** Consider operator $\mathcal{A}^{(5)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{28})^{(5)}$, $T_i^0 \le (\hat{Q}_{28})^{(5)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| By<br>$\bar{G}_{28}(t) = G_{28}^0 + \int_0^t \left[ (a_{28})^{(5)} G_{29}(s_{(28)}) - \left( (a'_{28})^{(5)} + a''_{28})^{(5)} (T_{29}(s_{(28)}), s_{(28)}) \right) G_{28}(s_{(28)}) \right] ds_{(28)}$ | 163 |
| $\bar{G}_{29}(t) = G_{29}^0 + \int_0^t \left[ (a_{29})^{(5)} G_{28}(s_{(28)}) - \left( (a'_{29})^{(5)} + (a''_{29})^{(5)} (T_{29}(s_{(28)}), s_{(28)}) \right) G_{29}(s_{(28)}) \right] ds_{(28)}$ | |
| $\bar{G}_{30}(t) = G_{30}^0 + \int_0^t \left[ (a_{30})^{(5)} G_{29}(s_{(28)}) - \left( (a'_{30})^{(5)} + (a''_{30})^{(5)} (T_{29}(s_{(28)}), s_{(28)}) \right) G_{30}(s_{(28)}) \right] ds_{(28)}$ | |
| $\bar{T}_{28}(t) = T_{28}^0 + \int_0^t \left[ (b_{28})^{(5)} T_{29}(s_{(28)}) - \left( (b'_{28})^{(5)} - (b''_{28})^{(5)} (G_{31}(s_{(28)}), s_{(28)}) \right) T_{28}(s_{(28)}) \right] ds_{(28)}$ | |
| $\bar{T}_{29}(t) = T_{29}^0 + \int_0^t \left[ (b_{29})^{(5)} T_{28}(s_{(28)}) - \left( (b'_{29})^{(5)} - (b''_{29})^{(5)} (G_{31}(s_{(28)}), s_{(28)}) \right) T_{29}(s_{(28)}) \right] ds_{(28)}$ | |
| $\bar{T}_{30}(t) = T_{30}^0 + \int_0^t \left[ (b_{30})^{(5)} T_{29}(s_{(28)}) - \left( (b'_{30})^{(5)} - (b''_{30})^{(5)} (G_{31}(s_{(28)}), s_{(28)}) \right) T_{30}(s_{(28)}) \right] ds_{(28)}$<br>Where $s_{(28)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:** | |

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| Consider operator $\mathcal{A}^{(6)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{32})^{(6)}$ , $T_i^0 \leq (\hat{Q}_{32})^{(6)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ | |
| By<br>$\bar{G}_{32}(t) = G_{32}^0 + \int_0^t \left[ (a_{32})^{(6)} G_{33}(s_{(32)}) - \left( (a'_{32})^{(6)} + a''_{32})^{(6)} (T_{33}(s_{(32)}), s_{(32)}) \right) G_{32}(s_{(32)}) \right] ds_{(32)}$ | 164 |
| $\bar{G}_{33}(t) = G_{33}^0 + \int_0^t \left[ (a_{33})^{(6)} G_{32}(s_{(32)}) - \left( (a'_{33})^{(6)} + (a''_{33})^{(6)} (T_{33}(s_{(32)}), s_{(32)}) \right) G_{33}(s_{(32)}) \right] ds_{(32)}$ | |
| $\bar{G}_{34}(t) = G_{34}^0 + \int_0^t \left[ (a_{34})^{(6)} G_{33}(s_{(32)}) - \left( (a'_{34})^{(6)} + (a''_{34})^{(6)} (T_{33}(s_{(32)}), s_{(32)}) \right) G_{34}(s_{(32)}) \right] ds_{(32)}$ | |
| $\bar{T}_{32}(t) = T_{32}^0 + \int_0^t \left[ (b_{32})^{(6)} T_{33}(s_{(32)}) - \left( (b'_{32})^{(6)} - (b''_{32})^{(6)} (G_{35}(s_{(32)}), s_{(32)}) \right) T_{32}(s_{(32)}) \right] ds_{(32)}$ | |
| $\bar{T}_{33}(t) = T_{33}^0 + \int_0^t \left[ (b_{33})^{(6)} T_{32}(s_{(32)}) - \left( (b'_{33})^{(6)} - (b''_{33})^{(6)} (G_{35}(s_{(32)}), s_{(32)}) \right) T_{33}(s_{(32)}) \right] ds_{(32)}$ | |
| $\bar{T}_{34}(t) = T_{34}^0 + \int_0^t \left[ (b_{34})^{(6)} T_{33}(s_{(32)}) - \left( (b'_{34})^{(6)} - (b''_{34})^{(6)} (G_{35}(s_{(32)}), s_{(32)}) \right) T_{34}(s_{(32)}) \right] ds_{(32)}$<br>Where $s_{(32)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(7)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{36})^{(7)}$ , $T_i^0 \leq (\hat{Q}_{36})^{(7)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ | |
| By<br>$\bar{G}_{36}(t) = G_{36}^0 + \int_0^t \left[ (a_{36})^{(7)} G_{37}(s_{(36)}) - \left( (a'_{36})^{(7)} + a''_{36})^{(7)} (T_{37}(s_{(36)}), s_{(36)}) \right) G_{36}(s_{(36)}) \right] ds_{(36)}$ | 165 |
| $\bar{G}_{37}(t) = G_{37}^0 + \int_0^t \left[ (a_{37})^{(7)} G_{36}(s_{(36)}) - \left( (a'_{37})^{(7)} + (a''_{37})^{(7)} (T_{37}(s_{(36)}), s_{(36)}) \right) G_{37}(s_{(36)}) \right] ds_{(36)}$ | |
| $\bar{G}_{38}(t) = G_{38}^0 + \int_0^t \left[ (a_{38})^{(7)} G_{37}(s_{(36)}) - \left( (a'_{38})^{(7)} + (a''_{38})^{(7)} (T_{37}(s_{(36)}), s_{(36)}) \right) G_{38}(s_{(36)}) \right] ds_{(36)}$ | |

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| $\bar{T}_{36}(t) = T_{36}^{0} + \int_0^t \left[(b_{36})^{(7)} T_{37}(s_{(36)}) - \left((b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{36}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{37}(t) = T_{37}^{0} + \int_0^t \left[(b_{37})^{(7)} T_{36}(s_{(36)}) - \left((b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{37}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{38}(t) = T_{38}^{0} + \int_0^t \left[(b_{38})^{(7)} T_{37}(s_{(36)}) - \left((b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{38}(s_{(36)})\right] ds_{(36)}$<br>Where $s_{(36)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(8)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \le (\hat{P}_{40})^{(8)}$ , $T_i^0 \le (\hat{Q}_{40})^{(8)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{40})^{(8)} e^{(\hat{M}_{40})^{(8)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{40})^{(8)} e^{(\hat{M}_{40})^{(8)} t}$ | |
| By<br>$\bar{G}_{40}(t) = G_{40}^0 + \int_0^t \left[(a_{40})^{(8)} G_{41}(s_{(40)}) - \left((a_{40}^{'})^{(8)} + a_{40}^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\right) G_{40}(s_{(40)})\right] ds_{(40)}$ | 166 |
| $\bar{G}_{41}(t) = G_{41}^0 + \int_0^t \left[(a_{41})^{(8)} G_{40}(s_{(40)}) - \left((a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\right) G_{41}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{G}_{42}(t) = G_{42}^0 + \int_0^t \left[(a_{42})^{(8)} G_{41}(s_{(40)}) - \left((a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\right) G_{42}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{T}_{40}(t) = T_{40}^{0} + \int_0^t \left[(b_{40})^{(8)} T_{41}(s_{(40)}) - \left((b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}(G_{43}(s_{(40)}), s_{(40)})\right) T_{40}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{T}_{41}(t) = T_{41}^{0} + \int_0^t \left[(b_{41})^{(8)} T_{40}(s_{(40)}) - \left((b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}(G_{43}(s_{(40)}), s_{(40)})\right) T_{41}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{T}_{42}(t) = T_{42}^{0} + \int_0^t \left[(b_{42})^{(8)} T_{41}(s_{(40)}) - \left((b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}(G_{43}(s_{(40)}), s_{(40)})\right) T_{42}(s_{(40)})\right] ds_{(40)}$<br>Where $s_{(40)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(9)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 166 A |

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| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{44})^{(9)}$ , $T_i^0 \leq (\hat{Q}_{44})^{(9)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} t}$ | |
| By<br>$\bar{G}_{44}(t) = G_{44}^0 + \int_0^t \left[(a_{44})^{(9)} G_{45}(s_{(44)}) - \left((a'_{44})^{(9)} + a''_{44})^{(9)} (T_{45}(s_{(44)}), s_{(44)})\right) G_{44}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{G}_{45}(t) = G_{45}^0 + \int_0^t \left[(a_{45})^{(9)} G_{44}(s_{(44)}) - \left((a'_{45})^{(9)} + (a''_{45})^{(9)} (T_{45}(s_{(44)}), s_{(44)})\right) G_{45}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{G}_{46}(t) = G_{46}^0 + \int_0^t \left[(a_{46})^{(9)} G_{45}(s_{(44)}) - \left((a'_{46})^{(9)} + (a''_{46})^{(9)} (T_{45}(s_{(44)}), s_{(44)})\right) G_{46}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{44}(t) = T_{44}^0 + \int_0^t \left[(b_{44})^{(9)} T_{45}(s_{(44)}) - \left((b'_{44})^{(9)} - (b''_{44})^{(9)} (G_{47}(s_{(44)}), s_{(44)})\right) T_{44}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{45}(t) = T_{45}^0 + \int_0^t \left[(b_{45})^{(9)} T_{44}(s_{(44)}) - \left((b'_{45})^{(9)} - (b''_{45})^{(9)} (G_{47}(s_{(44)}), s_{(44)})\right) T_{45}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{46}(t) = T_{46}^0 + \int_0^t \left[(b_{46})^{(9)} T_{45}(s_{(44)}) - \left((b'_{46})^{(9)} - (b''_{46})^{(9)} (G_{47}(s_{(44)}), s_{(44)})\right) T_{46}(s_{(44)})\right] ds_{(44)}$<br>Where $s_{(44)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| The operator $\mathcal{A}^{(1)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{13}(t) \leq G_{13}^0 + \int_0^t \left[(a_{13})^{(1)} \left(G_{14}^0 + (\hat{P}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} s_{(13)}}\right)\right] ds_{(13)} =$<br>$\left(1 + (a_{13})^{(1)} t\right) G_{14}^0 + \frac{(a_{13})^{(1)} (\hat{P}_{13})^{(1)}}{(\hat{M}_{13})^{(1)}} \left(e^{(\hat{M}_{13})^{(1)} t} - 1\right)$ | 167 |
| From which it follows that<br>$(G_{13}(t) - G_{13}^0) e^{-(\hat{M}_{13})^{(1)} t} \leq \frac{(a_{13})^{(1)}}{(\hat{M}_{13})^{(1)}} \left[\left((\hat{P}_{13})^{(1)} + G_{14}^0\right) e^{\left(-\frac{(\hat{P}_{13})^{(1)} + G_{14}^0}{G_{14}^0}\right)} + (\hat{P}_{13})^{(1)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 1 | 168 |
| Analogous inequalities hold also for $G_{14}$ , $G_{15}, T_{13}, T_{14}, T_{15}$ | |
| The operator $\mathcal{A}^{(2)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that | |
| $G_{16}(t) \leq G_{16}^0 + \int_0^t \left[(a_{16})^{(2)} \left(G_{17}^0 + (\hat{P}_{16})^{(6)} e^{(\hat{M}_{16})^{(2)} s_{(16)}}\right)\right] ds_{(16)} =$<br>$\left(1 + (a_{16})^{(2)} t\right) G_{17}^0 + \frac{(a_{16})^{(2)} (\hat{P}_{16})^{(2)}}{(\hat{M}_{16})^{(2)}} \left(e^{(\hat{M}_{16})^{(2)} t} - 1\right)$ | 169 |
| From which it follows that | 170 |

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| $(G_{16}(t)-G_{16}^0)e^{-(\bar{M}_{16})^{(2)}t} \le \frac{(a_{16})^{(2)}}{(\bar{M}_{16})^{(2)}}\left[((\hat{P}_{16})^{(2)}+G_{17}^0)e^{\left(-\frac{(\hat{P}_{16})^{(2)}+G_{17}^0}{G_{17}^0}\right)}+(\hat{P}_{16})^{(2)}\right]$ | |
| Analogous inequalities hold also for $G_{17}, G_{18}, T_{16}, T_{17}, T_{18}$ | |
| The operator $\mathcal{A}^{(3)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{20}(t) \le G_{20}^0 + \int_0^t \left[(a_{20})^{(3)}\left(G_{21}^0+(\hat{P}_{20})^{(3)}e^{(\bar{M}_{20})^{(3)}s_{(20)}}\right)\right] ds_{(20)} =$<br>$(1+(a_{20})^{(3)}t)G_{21}^0 + \frac{(a_{20})^{(3)}(\hat{P}_{20})^{(3)}}{(\bar{M}_{20})^{(3)}}\left(e^{(\bar{M}_{20})^{(3)}t}-1\right)$ | 171 |
| From which it follows that<br>$(G_{20}(t)-G_{20}^0)e^{-(\bar{M}_{20})^{(3)}t} \le \frac{(a_{20})^{(3)}}{(\bar{M}_{20})^{(3)}}\left[((\hat{P}_{20})^{(3)}+G_{21}^0)e^{\left(-\frac{(\hat{P}_{20})^{(3)}+G_{21}^0}{G_{21}^0}\right)}+(\hat{P}_{20})^{(3)}\right]$ | 172 |
| Analogous inequalities hold also for $G_{21}, G_{22}, T_{20}, T_{21}, T_{22}$ | |
| The operator $\mathcal{A}^{(4)}$ maps the space of functions satisfying into itself .Indeed it is obvious that<br>$G_{24}(t) \le G_{24}^0 + \int_0^t \left[(a_{24})^{(4)}\left(G_{25}^0+(\hat{P}_{24})^{(4)}e^{(\bar{M}_{24})^{(4)}s_{(24)}}\right)\right] ds_{(24)} =$<br>$(1+(a_{24})^{(4)}t)G_{25}^0 + \frac{(a_{24})^{(4)}(\hat{P}_{24})^{(4)}}{(\bar{M}_{24})^{(4)}}\left(e^{(\bar{M}_{24})^{(4)}t}-1\right)$ | 173 |
| From which it follows that<br>$(G_{24}(t)-G_{24}^0)e^{-(\bar{M}_{24})^{(4)}t} \le \frac{(a_{24})^{(4)}}{(\bar{M}_{24})^{(4)}}\left[((\hat{P}_{24})^{(4)}+G_{25}^0)e^{\left(-\frac{(\hat{P}_{24})^{(4)}+G_{25}^0}{G_{25}^0}\right)}+(\hat{P}_{24})^{(4)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 4 | 174 |
| The operator $\mathcal{A}^{(5)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{28}(t) \le G_{28}^0 + \int_0^t \left[(a_{28})^{(5)}\left(G_{29}^0+(\hat{P}_{28})^{(5)}e^{(\bar{M}_{28})^{(5)}s_{(28)}}\right)\right] ds_{(28)} =$<br>$(1+(a_{28})^{(5)}t)G_{29}^0 + \frac{(a_{28})^{(5)}(\hat{P}_{28})^{(5)}}{(\bar{M}_{28})^{(5)}}\left(e^{(\bar{M}_{28})^{(5)}t}-1\right)$ | |
| From which it follows that<br>$(G_{28}(t)-G_{28}^0)e^{-(\bar{M}_{28})^{(5)}t} \le \frac{(a_{28})^{(5)}}{(\bar{M}_{28})^{(5)}}\left[((\hat{P}_{28})^{(5)}+G_{29}^0)e^{\left(-\frac{(\hat{P}_{28})^{(5)}+G_{29}^0}{G_{29}^0}\right)}+(\hat{P}_{28})^{(5)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 5 | 175 |
| The operator $\mathcal{A}^{(6)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious | 176 |

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| that<br>$G_{32}(t) \le G_{32}^0 + \int_0^t \left[(a_{32})^{(6)}\left(G_{33}^0 + (\hat{P}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} s_{(32)}}\right)\right] ds_{(32)} =$<br>$\left(1 + (a_{32})^{(6)} t\right) G_{33}^0 + \frac{(a_{32})^{(6)} (\hat{P}_{32})^{(6)}}{(\hat{M}_{32})^{(6)}}\left(e^{(\hat{M}_{32})^{(6)} t} - 1\right)$ | |
| From which it follows that<br>$(G_{32}(t) - G_{32}^0) e^{-(\hat{M}_{32})^{(6)} t} \le \frac{(a_{32})^{(6)}}{(\hat{M}_{32})^{(6)}}\left[\left((\hat{P}_{32})^{(6)} + G_{33}^0\right) e^{\left(-\frac{(\hat{P}_{32})^{(6)} + G_{33}^0}{G_{33}^0}\right)} + (\hat{P}_{32})^{(6)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 6<br>Analogous inequalities hold also for $G_{25}, G_{26}, T_{24}, T_{25}, T_{26}$ | 177 |
| (b) The operator $\mathcal{A}^{(7)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{36}(t) \le G_{36}^0 + \int_0^t \left[(a_{36})^{(7)}\left(G_{37}^0 + (\hat{P}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} s_{(36)}}\right)\right] ds_{(36)} =$<br>$\left(1 + (a_{36})^{(7)} t\right) G_{37}^0 + \frac{(a_{36})^{(7)} (\hat{P}_{36})^{(7)}}{(\hat{M}_{36})^{(7)}}\left(e^{(\hat{M}_{36})^{(7)} t} - 1\right)$ | 178 |
| From which it follows that<br>$(G_{36}(t) - G_{36}^0) e^{-(\hat{M}_{36})^{(7)} t} \le \frac{(a_{36})^{(7)}}{(\hat{M}_{36})^{(7)}}\left[\left((\hat{P}_{36})^{(7)} + G_{37}^0\right) e^{\left(-\frac{(\hat{P}_{36})^{(7)} + G_{37}^0}{G_{37}^0}\right)} + (\hat{P}_{36})^{(7)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 7 | |
| The operator $\mathcal{A}^{(8)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that | |
| $G_{40}(t) \le G_{40}^0 + \int_0^t \left[(a_{40})^{(8)}\left(G_{41}^0 + (\hat{P}_{40})^{(8)} e^{(\hat{M}_{40})^{(8)} s_{(40)}}\right)\right] ds_{(40)} =$<br>$\left(1 + (a_{40})^{(8)} t\right) G_{41}^0 + \frac{(a_{40})^{(8)} (\hat{P}_{40})^{(8)}}{(\hat{M}_{40})^{(8)}}\left(e^{(\hat{M}_{40})^{(8)} t} - 1\right)$ | 180 |
| From which it follows that<br>$(G_{40}(t) - G_{40}^0) e^{-(\hat{M}_{40})^{(8)} t} \le \frac{(a_{40})^{(8)}}{(\hat{M}_{40})^{(8)}}\left[\left((\hat{P}_{40})^{(8)} + G_{41}^0\right) e^{\left(-\frac{(\hat{P}_{40})^{(8)} + G_{41}^0}{G_{41}^0}\right)} + (\hat{P}_{40})^{(8)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 8<br>Analogous inequalities hold also for $G_{41}, G_{42}, T_{40}, T_{41}, T_{42}$ | 181 |
| The operator $\mathcal{A}^{(9)}$ maps the space of functions satisfying 34,35,36 into itself .Indeed it is obvious that<br>$G_{44}(t) \le G_{44}^0 + \int_0^t \left[(a_{44})^{(9)}\left(G_{45}^0 + (\hat{P}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} s_{(44)}}\right)\right] ds_{(44)} =$ | |

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| $\left(1+(a_{44})^{(9)}t\right)G_{45}^{0}+\frac{(a_{44})^{(9)}(\hat{P}_{44})^{(9)}}{(\hat{M}_{44})^{(9)}}\left(e^{(\hat{M}_{44})^{(9)}t}-1\right)$ | |
| From which it follows that<br>$(G_{44}(t)-G_{44}^{0})e^{-(\hat{M}_{44})^{(9)}t}\leq\frac{(a_{44})^{(9)}}{(\hat{M}_{44})^{(9)}}\left[\left((\hat{P}_{44})^{(9)}+G_{45}^{0}\right)e^{\left(-\frac{(\hat{P}_{44})^{(9)}+G_{45}^{0}}{G_{45}^{0}}\right)}+(\hat{P}_{44})^{(9)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 9<br>Analogous inequalities hold also for $G_{45}, G_{46}, T_{44}, T_{45}, T_{46}$ | |
| It is now sufficient to take $\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}}<1$ and to choose<br>$(\hat{P}_{13})^{(1)}$ and $(\hat{Q}_{13})^{(1)}$ large to have | 182 |
| $\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}\left[(\hat{P}_{13})^{(1)}+\left((\hat{P}_{13})^{(1)}+G_j^0\right)e^{-\left(\frac{(\hat{P}_{13})^{(1)}+G_j^0}{G_j^0}\right)}\right]\leq(\hat{P}_{13})^{(1)}$ | 183 |
| $\frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}}\left[\left((\hat{Q}_{13})^{(1)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{13})^{(1)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{13})^{(1)}\right]\leq(\hat{Q}_{13})^{(1)}$ | 184 |
| In order that the operator $\mathcal{A}^{(1)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(1)}$ is a contraction with respect to the metric<br>$d\left(\left(G^{(1)},T^{(1)}\right),\left(G^{(2)},T^{(2)}\right)\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\hat{M}_{13})^{(1)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\hat{M}_{13})^{(1)}t}\}$ | 185 |
| Indeed if we denote<br>**Definition of** $\tilde{G},\tilde{T}:\left(\tilde{G},\tilde{T}\right)=\mathcal{A}^{(1)}(G,T)$<br>It results<br>$\left\|\tilde{G}_{13}^{(1)}-\tilde{G}_i^{(2)}\right\|\leq\int_0^t(a_{13})^{(1)}\left\|G_{14}^{(1)}-G_{14}^{(2)}\right\|e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{(\hat{M}_{13})^{(1)}s_{(13)}}ds_{(13)}+$<br>$\int_0^t\{(a_{13}')^{(1)}\left\|G_{13}^{(1)}-G_{13}^{(2)}\right\|e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{-(\hat{M}_{13})^{(1)}s_{(13)}}+$<br>$(a_{13}'')^{(1)}\left(T_{14}^{(1)},s_{(13)}\right)\left\|G_{13}^{(1)}-G_{13}^{(2)}\right\|e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{(\hat{M}_{13})^{(1)}s_{(13)}}+$<br>$G_{13}^{(2)}\left\\|(a_{13}'')^{(1)}\left(T_{14}^{(1)},s_{(13)}\right)-(a_{13}'')^{(1)}\left(T_{14}^{(2)},s_{(13)}\right)\right\\|\ e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{(\hat{M}_{13})^{(1)}s_{(13)}}\}ds_{(13)}$ | |

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| Where $s_{(13)}$ represents integrand that is integrated over the interval $[0, t]$<br><br>From the hypotheses it follows | |
| $\left\|G^{(1)} - G^{(2)}\right\| e^{-(\widehat{M}_{13})^{(1)}t} \leq$<br>$\frac{1}{(\widehat{M}_{13})^{(1)}}\left((a_{13})^{(1)} + (a'_{13})^{(1)} + (\widehat{A}_{13})^{(1)} + (\widehat{P}_{13})^{(1)}(\widehat{k}_{13})^{(1)}\right) d\left(\left(G^{(1)}, T^{(1)};\ G^{(2)}, T^{(2)}\right)\right)$<br><br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 186 |
| **Remark 1:** The fact that we supposed $(a''_{13})^{(1)}$ and $(b''_{13})^{(1)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{13})^{(1)} e^{(\widehat{M}_{13})^{(1)}t}$ *and* $(\widehat{Q}_{13})^{(1)} e^{(\widehat{M}_{13})^{(1)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a''_i)^{(1)}$ and $(b''_i)^{(1)}, i = 13,14,15$ depend only on $T_{14}$ and respectively on $G(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 2:** There does not exist any $t$ where $G_i(t) = 0$ *and* $T_i(t) = 0$<br><br>From 19 to 24 it results<br><br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a'_i)^{(1)} - (a''_i)^{(1)}(T_{14}(s_{(13)}), s_{(13)})\} ds_{(13)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{\left(-(b'_i)^{(1)}t\right)} > 0$ for t > 0 | |
| **Definition of** $\left((\widehat{M}_{13})^{(1)}\right)_1, \left((\widehat{M}_{13})^{(1)}\right)_2$ *and* $\left((\widehat{M}_{13})^{(1)}\right)_3$ :<br><br>**Remark 3**: if $G_{13}$ is bounded, the same property have also $G_{14}$ *and* $G_{15}$ . indeed if<br><br>$G_{13} < (\widehat{M}_{13})^{(1)}$ it follows $\frac{dG_{14}}{dt} \leq \left((\widehat{M}_{13})^{(1)}\right)_1 - (a'_{14})^{(1)} G_{14}$ and by integrating<br><br>$G_{14} \leq \left((\widehat{M}_{13})^{(1)}\right)_2 = G_{14}^0 + 2(a_{14})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_1 / (a'_{14})^{(1)}$<br><br>In the same way , one can obtain<br><br>$G_{15} \leq \left((\widehat{M}_{13})^{(1)}\right)_3 = G_{15}^0 + 2(a_{15})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_2 / (a'_{15})^{(1)}$<br><br>If $G_{14}$ *or* $G_{15}$ is bounded, the same property follows for $G_{13}$ , $G_{15}$ and $G_{13}$ , $G_{14}$ respectively. | 187 |
| **Remark 4:** If $G_{13}$ *is* bounded, from below, the same property holds for $G_{14}$ *and* $G_{15}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{14}$ is bounded from below. | 188 |
| **Remark 5:** If $T_{13}$ is bounded from below and $\lim_{t\to\infty}((b''_i)^{(1)}(G(t), t)) = (b'_{14})^{(1)}$ then $T_{14} \to \infty$.<br><br>**Definition of** $(m)^{(1)}$ and $\varepsilon_1$ : | 189 |

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| Indeed let $t_1$ be so that for $t > t_1$<br>$(b_{14})^{(1)} - (b_i'')^{(1)}(G(t),t) < \varepsilon_1, T_{13}(t) > (m)^{(1)}$ | |
| Then $\frac{dT_{14}}{dt} \geq (a_{14})^{(1)}(m)^{(1)} - \varepsilon_1 T_{14}$ which leads to<br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{\varepsilon_1}\right)(1 - e^{-\varepsilon_1 t}) + T_{14}^0 e^{-\varepsilon_1 t}$ If we take t such that $e^{-\varepsilon_1 t} = \frac{1}{2}$ it results<br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{2}\right), \quad t = log\frac{2}{\varepsilon_1}$ By taking now $\varepsilon_1$ sufficiently small one sees that $T_{14}$ is unbounded.<br>The same property holds for $T_{15}$ if $\lim_{t\to\infty}(b_{15}'')^{(1)}(G(t),t) = (b_{15}')^{(1)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |
| It is now sufficient to take $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}, \frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ and to choose<br>$(\hat{P}_{16})^{(2)}$ $and$ $(\hat{Q}_{16})^{(2)}$ large to have | 190 |
| $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[(\hat{P}_{16})^{(2)} + \left((\hat{P}_{16})^{(2)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{16})^{(2)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{16})^{(2)}$ | 191 |
| $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[\left((\hat{Q}_{16})^{(2)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{16})^{(2)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{16})^{(2)}\right] \leq (\hat{Q}_{16})^{(2)}$ | 192 |
| In order that the operator $\mathcal{A}^{(2)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | 193 |
| The operator $\mathcal{A}^{(2)}$ is a contraction with respect to the metric<br>$d\left(((G_{19})^{(1)},(T_{19})^{(1)}),((G_{19})^{(2)},(T_{19})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\hat{M}_{16})^{(2)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\hat{M}_{16})^{(2)}t}\}$ | 194 |
| Indeed if we denote<br>**Definition of** $\widetilde{G_{19}}, \widetilde{T_{19}}$ : $\left(\widetilde{G_{19}}, \widetilde{T_{19}}\right) = \mathcal{A}^{(2)}(G_{19}, T_{19})$ | 195 |

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| It results<br>$\left\|\tilde{G}_{16}^{(1)} - \tilde{G}_{i}^{(2)}\right\| \leq \int_0^t (a_{16})^{(2)} \left\|G_{17}^{(1)} - G_{17}^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{(\widehat{M}_{16})^{(2)} s_{(16)}}\, ds_{(16)} +$<br>$\int_0^t \{(a_{16}^{'})^{(2)} \left\|G_{16}^{(1)} - G_{16}^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} +$<br>$(a_{16}^{''})^{(2)} \left(T_{17}^{(1)}, s_{(16)}\right) \left\|G_{16}^{(1)} - G_{16}^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{(\widehat{M}_{16})^{(2)} s_{(16)}} +$<br>$G_{16}^{(2)} \left\|(a_{16}^{''})^{(2)} \left(T_{17}^{(1)}, s_{(16)}\right) - (a_{16}^{''})^{(2)} \left(T_{17}^{(2)}, s_{(16)}\right)\right\| \; e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{(\widehat{M}_{16})^{(2)} s_{(16)}} \} ds_{(16)}$ | 196 |
| Where $s_{(16)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | 197 |
| $\left\|(G_{19})^{(1)} - (G_{19})^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} t} \leq$<br>$\frac{1}{(\widehat{M}_{16})^{(2)}} \left((a_{16})^{(2)} + (a_{16}^{'})^{(2)} + (\widehat{A}_{16})^{(2)} + (\widehat{P}_{16})^{(2)} (\widehat{k}_{16})^{(2)}\right) d\left(\left((G_{19})^{(1)}, (T_{19})^{(1)}; (G_{19})^{(2)}, (T_{19})^{(2)}\right)\right)$ | |
| And analogous inequalities for $G_i$ and $T_i$. Taking into account the hypothesis the result follows | 198 |
| **Remark 6:** The fact that we supposed $(a_{16}^{''})^{(2)}$ and $(b_{16}^{''})^{(2)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{16})^{(2)} e^{(\widehat{M}_{16})^{(2)} t}$ and $(\widehat{Q}_{16})^{(2)} e^{(\widehat{M}_{16})^{(2)} t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}, i = 16,17,18$ depend only on $T_{17}$ and respectively on $(G_{19})$(and not on t) and hypothesis can replaced by a usual Lipschitz condition. | 199 |
| **Remark 7:** There does not exist any t where $G_i(t) = 0$ and $T_i(t) = 0$<br>it results<br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(2)} - (a_i^{''})^{(2)} (T_{17}(s_{(16)}), s_{(16)})\} ds_{(16)}\right]} \geq 0$<br>$T_i(t) \geq T_i^0 e^{(-(b_i^{'})^{(2)} t)} > 0 \quad \text{for } t > 0$ | 200 |
| **Definition of** $\left((\widehat{M}_{16})^{(2)}\right)_1$, $\left((\widehat{M}_{16})^{(2)}\right)_2$ and $\left((\widehat{M}_{16})^{(2)}\right)_3$ :<br>**Remark 8:** if $G_{16}$ is bounded, the same property have also $G_{17}$ and $G_{18}$ . indeed if<br>$G_{16} < (\widehat{M}_{16})^{(2)}$ it follows $\frac{dG_{17}}{dt} \leq \left((\widehat{M}_{16})^{(2)}\right)_1 - (a_{17}^{'})^{(2)} G_{17}$ and by integrating<br>$G_{17} \leq \left((\widehat{M}_{16})^{(2)}\right)_2 = G_{17}^0 + 2(a_{17})^{(2)} \left((\widehat{M}_{16})^{(2)}\right)_1 / (a_{17}^{'})^{(2)}$<br>In the same way , one can obtain | 201 |

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| $G_{18} \leq \left((\widehat{M}_{16})^{(2)}\right)_3 = G_{18}^0 + 2(a_{18})^{(2)}\left((\widehat{M}_{16})^{(2)}\right)_2/(a_{18}')^{(2)}$<br><br>If $G_{17}$ or $G_{18}$ is bounded, the same property follows for $G_{16}$, $G_{18}$ and $G_{16}$, $G_{17}$ respectively. | |
| **Remark 9**: If $G_{16}$ is bounded, from below, the same property holds for $G_{17}$ and $G_{18}$. The proof is analogous with the preceding one. An analogous property is true if $G_{17}$ is bounded from below. | 202 |
| **Remark 10:** If $T_{16}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(2)}((G_{19})(t),t)) = (b_{17}')^{(2)}$ then $T_{17} \to \infty$.<br><br>**Definition of** $(m)^{(2)}$ and $\varepsilon_2$ :<br><br>Indeed let $t_2$ be so that for $t > t_2$<br><br>$(b_{17})^{(2)} - (b_i'')^{(2)}((G_{19})(t),t) < \varepsilon_2, T_{16}(t) > (m)^{(2)}$ | 203 |
| Then $\frac{dT_{17}}{dt} \geq (a_{17})^{(2)}(m)^{(2)} - \varepsilon_2 T_{17}$ which leads to<br><br>$T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{\varepsilon_2}\right)(1 - e^{-\varepsilon_2 t}) + T_{17}^0 e^{-\varepsilon_2 t}$ If we take t such that $e^{-\varepsilon_2 t} = \frac{1}{2}$ it results | 204 |
| $T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{2}\right)$, $t = \log\frac{2}{\varepsilon_2}$ By taking now $\varepsilon_2$ sufficiently small one sees that $T_{17}$ is unbounded.<br>The same property holds for $T_{18}$ if $\lim_{t\to\infty}(b_{18}'')^{(2)}\left((G_{19})(t),t\right) = (b_{18}')^{(2)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 205 |
| It is now sufficient to take $\frac{(a_i)^{(3)}}{(\bar{M}_{20})^{(3)}}$, $\frac{(b_i)^{(3)}}{(\bar{M}_{20})^{(3)}} < 1$ and to choose<br><br>$(\hat{P}_{20})^{(3)}$ and $(\hat{Q}_{20})^{(3)}$ large to have | 207 |
| $\frac{(a_i)^{(3)}}{(\bar{M}_{20})^{(3)}}\left[(\hat{P}_{20})^{(3)} + \left((\hat{P}_{20})^{(3)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{20})^{(3)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{20})^{(3)}$ | 208 |
| $\frac{(b_i)^{(3)}}{(\bar{M}_{20})^{(3)}}\left[\left((\hat{Q}_{20})^{(3)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{20})^{(3)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{20})^{(3)}\right] \leq (\hat{Q}_{20})^{(3)}$ | 209 |
| In order that the operator $\mathcal{A}^{(3)}$ transforms the space of sextuples of functions $G_i$, $T_i$ satisfying Equations into itself | 210 |
| The operator $\mathcal{A}^{(3)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{23})^{(1)},(T_{23})^{(1)}),((G_{23})^{(2)},(T_{23})^{(2)})\right) =$<br><br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t) - G_i^{(2)}(t)\right|e^{-(\bar{M}_{20})^{(3)}t}, \max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t) - T_i^{(2)}(t)\right|e^{-(\bar{M}_{20})^{(3)}t}\}$ | 211 |

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| Indeed if we denote<br><br>**Definition of** $\widetilde{G_{23}}, \widetilde{T_{23}} :\left( \widetilde{(G_{23})}, \widetilde{(T_{23})} \right) = \mathcal{A}^{(3)}\left((G_{23}), (T_{23})\right)$ | 212 |
| It results<br><br>$\left\|\tilde{G}_{20}^{(1)} - \tilde{G}_{i}^{(2)}\right\| \leq \int_0^t (a_{20})^{(3)} \left\|G_{21}^{(1)} - G_{21}^{(2)}\right\| e^{-(\widehat{M}_{20})^{(3)} s_{(20)}} e^{(\widehat{M}_{20})^{(3)} s_{(20)}}\, ds_{(20)} +$<br><br>$\int_0^t \{(a_{20}^{'})^{(3)} \left\|G_{20}^{(1)} - G_{20}^{(2)}\right\| e^{-(\widehat{M}_{20})^{(3)} s_{(20)}} e^{-(\widehat{M}_{20})^{(3)} s_{(20)}} +$<br><br>$(a_{20}^{''})^{(3)}\left(T_{21}^{(1)}, s_{(20)}\right)\left\|G_{20}^{(1)} - G_{20}^{(2)}\right\| e^{-(\widehat{M}_{20})^{(3)} s_{(20)}} e^{(\widehat{M}_{20})^{(3)} s_{(20)}} +$<br><br>$G_{20}^{(2)} \left\|(a_{20}^{''})^{(3)}\left(T_{21}^{(1)}, s_{(20)}\right) - (a_{20}^{''})^{(3)}\left(T_{21}^{(2)}, s_{(20)}\right)\right\| \; e^{-(\widehat{M}_{20})^{(3)} s_{(20)}} e^{(\widehat{M}_{20})^{(3)} s_{(20)}}\} ds_{(20)}$<br><br>Where $s_{(20)}$ represents integrand that is integrated over the interval $[0, t]$<br><br>From the hypotheses it follows | 213 |
| $\left\|G_{23}{}^{(1)} - G_{23}{}^{(2)}\right\| e^{-(\widehat{M}_{20})^{(3)} t} \leq$<br>$\frac{1}{(\widehat{M}_{20})^{(3)}}\left((a_{20})^{(3)} + (a_{20}^{'})^{(3)} + (\widehat{A}_{20})^{(3)} + (\widehat{P}_{20})^{(3)} (\widehat{k}_{20})^{(3)}\right) d\left(\left((G_{23})^{(1)}, (T_{23})^{(1)};\ (G_{23})^{(2)}, (T_{23})^{(2)}\right)\right)$<br><br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 214 |
| **Remark 11:** The fact that we supposed $(a_{20}^{''})^{(3)}$ and $(b_{20}^{''})^{(3)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{20})^{(3)} e^{(\widehat{M}_{20})^{(3)} t}$ *and* $(\widehat{Q}_{20})^{(3)} e^{(\widehat{M}_{20})^{(3)} t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(3)}$ and $(b_i^{''})^{(3)}, i = 20,21,22$ depend only on $T_{21}$ and respectively on $(G_{23})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 215 |
| **Remark 12:** There does not exist any $t$ where $G_i\,(t) = 0$ *and* $T_i\,(t) = 0$<br><br>it results<br><br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(3)} - (a_i^{''})^{(3)}(T_{21}(s_{(20)}), s_{(20)})\} ds_{(20)}\right]} \geq 0$<br><br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(3)} t\right)} > 0$ for t > 0 | 216 |
| **Definition of** $\left((\widehat{M}_{20})^{(3)}\right)_1, \left((\widehat{M}_{20})^{(3)}\right)_2$ *and* $\left((\widehat{M}_{20})^{(3)}\right)_3$ :<br><br>**Remark 13:** if $G_{20}$ is bounded, the same property have also $G_{21}$ *and* $G_{22}$ . indeed if<br><br>$G_{20} < (\widehat{M}_{20})^{(3)}$ it follows $\frac{dG_{21}}{dt} \leq \left((\widehat{M}_{20})^{(3)}\right)_1 - (a_{21}^{'})^{(3)} G_{21}$ and by integrating | 217 |

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| $G_{21} \leq \left((\widehat{M}_{20})^{(3)}\right)_2 = G_{21}^0 + 2(a_{21})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_1/(a_{21}')^{(3)}$<br><br>In the same way , one can obtain<br><br>$G_{22} \leq \left((\widehat{M}_{20})^{(3)}\right)_3 = G_{22}^0 + 2(a_{22})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_2/(a_{22}')^{(3)}$<br><br>If $G_{21}$ *or* $G_{22}$ is bounded, the same property follows for $G_{20}$ , $G_{22}$ and $G_{20}$ , $G_{21}$ respectively. | |
| **Remark 14:** If $G_{20}$ *is* bounded, from below, the same property holds for $G_{21}$ *and* $G_{22}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{21}$ is bounded from below. | 218 |
| **Remark 15:** If $T_{20}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(3)}\left((G_{23})(t),t\right) = (b_{21}^{'})^{(3)}$ then $T_{21} \to \infty$.<br><br>**Definition of** $(m)^{(3)}$ and $\varepsilon_3$ :<br><br>Indeed let $t_3$ be so that for $t > t_3$<br><br>$(b_{21})^{(3)} - (b_i^{''})^{(3)}\left((G_{23})(t),t\right) < \varepsilon_3, T_{20}(t) > (m)^{(3)}$ | 219 |
| Then $\frac{dT_{21}}{dt} \geq (a_{21})^{(3)}(m)^{(3)} - \varepsilon_3 T_{21}$ which leads to<br><br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{\varepsilon_3}\right)(1 - e^{-\varepsilon_3 t}) + T_{21}^0 e^{-\varepsilon_3 t}$ If we take t such that $e^{-\varepsilon_3 t} = \frac{1}{2}$ it results<br><br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{2}\right)$, $t = log\frac{2}{\varepsilon_3}$ By taking now $\varepsilon_3$ sufficiently small one sees that $T_{21}$ is unbounded.<br>The same property holds for $T_{22}$ if $\lim_{t\to\infty}(b_{22}^{''})^{(3)}\left((G_{23})(t),t\right) = (b_{22}^{'})^{(3)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 220 |
| It is now sufficient to take $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}$ , $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}} < 1$ and to choose<br><br>$(\widehat{P}_{24})^{(4)}$ and $(\widehat{Q}_{24})^{(4)}$ large to have | 221 |
| $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[(\widehat{P}_{24})^{(4)} + \left((\widehat{P}_{24})^{(4)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{24})^{(4)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{24})^{(4)}$ | 222 |
| $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[\left((\widehat{Q}_{24})^{(4)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{24})^{(4)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{24})^{(4)}\right] \leq (\widehat{Q}_{24})^{(4)}$ | 223 |
| In order that the operator $\mathcal{A}^{(4)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 224 |
| The operator $\mathcal{A}^{(4)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{27})^{(1)},(T_{27})^{(1)}),((G_{27})^{(2)},(T_{27})^{(2)})\right) =$ | 225 |

$$\sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t)-G_i^{(2)}(t)\right|e^{-(\widehat{M}_{24})^{(4)}t},\max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t)-T_i^{(2)}(t)\right|e^{-(\widehat{M}_{24})^{(4)}t}\}$$

Indeed if we denote

**Definition of** $\widetilde{(G_{27})},\widetilde{(T_{27})}$ : $\left(\widetilde{(G_{27})},\widetilde{(T_{27})}\right)=\mathcal{A}^{(4)}((G_{27}),(T_{27}))$

It results

$$\left|\tilde{G}_{24}^{(1)}-\tilde{G}_i^{(2)}\right|\le\int_0^t(a_{24})^{(4)}\left|G_{25}^{(1)}-G_{25}^{(2)}\right|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}\,ds_{(24)}+$$

$$\int_0^t\{(a_{24}^{'})^{(4)}\left|G_{24}^{(1)}-G_{24}^{(2)}\right|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}+$$

$$(a_{24}^{''})^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)\left|G_{24}^{(1)}-G_{24}^{(2)}\right|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}+$$

$$G_{24}^{(2)}\left|(a_{24}^{''})^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)-(a_{24}^{''})^{(4)}\left(T_{25}^{(2)},s_{(24)}\right)\right|\ e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}\}ds_{(24)}$$

Where $s_{(24)}$ represents integrand that is integrated over the interval $[0,t]$

From the hypotheses on Equations it follows

$$\left|(G_{27})^{(1)}-(G_{27})^{(2)}\right|e^{-(\widehat{M}_{24})^{(4)}t}\le$$
$$\frac{1}{(\widehat{M}_{24})^{(4)}}\left((a_{24})^{(4)}+(a_{24}^{'})^{(4)}+(\widehat{A}_{24})^{(4)}+(\widehat{P}_{24})^{(4)}(\widehat{k}_{24})^{(4)}\right)d\left(((G_{27})^{(1)},(T_{27})^{(1)};\ (G_{27})^{(2)},(T_{27})^{(2)})\right)$$

And analogous inequalities for $G_i\ and\ T_i$. Taking into account the hypothesis the result follows

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**Remark 16:** The fact that we supposed $(a_{24}^{''})^{(4)}$ and $(b_{24}^{''})^{(4)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}\ and\ (\widehat{Q}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}$ respectively of $\mathbb{R}_+$.

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(4)}$ and $(b_i^{''})^{(4)}, i=24,25,26$ depend only on $T_{25}$ and respectively on $(G_{27})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition.

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**Remark 17:** There does not exist any $t$ where $G_i\ (t)=0\ and\ T_i\ (t)=0$

it results

$$G_i\ (t)\ge G_i^0e^{\left[-\int_0^t\{(a_i^{'})^{(4)}-(a_i^{''})^{(4)}(T_{25}(s_{(24)}),s_{(24)})\}ds_{(24)}\right]}\ge 0$$

$$T_i\ (t)\ge T_i^0e^{\left(-(b_i^{'})^{(4)}t\right)}>0\quad \text{for t} > 0$$

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**Definition of** $\left((\widehat{M}_{24})^{(4)}\right)_1,\left((\widehat{M}_{24})^{(4)}\right)_2\ and\ \left((\widehat{M}_{24})^{(4)}\right)_3$ :

**Remark 18:** if $G_{24}$ is bounded, the same property have also $G_{25}\ and\ G_{26}$ . indeed if

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| $G_{24} < (\widehat{M}_{24})^{(4)}$ it follows $\frac{dG_{25}}{dt} \leq \left((\widehat{M}_{24})^{(4)}\right)_1 - (a'_{25})^{(4)} G_{25}$ and by integrating<br>$G_{25} \leq \left((\widehat{M}_{24})^{(4)}\right)_2 = G_{25}^0 + 2(a_{25})^{(4)}\left((\widehat{M}_{24})^{(4)}\right)_1/(a'_{25})^{(4)}$<br>In the same way , one can obtain<br>$G_{26} \leq \left((\widehat{M}_{24})^{(4)}\right)_3 = G_{26}^0 + 2(a_{26})^{(4)}\left((\widehat{M}_{24})^{(4)}\right)_2/(a'_{26})^{(4)}$<br>If $G_{25}$ *or* $G_{26}$ is bounded, the same property follows for $G_{24}$ , $G_{26}$ and $G_{24}$ , $G_{25}$ respectively. | |
| **Remark 19:** If $G_{24}$ *is* bounded, from below, the same property holds for $G_{25}$ *and* $G_{26}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{25}$ is bounded from below. | 230 |
| **Remark 20:** If $T_{24}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(4)}((G_{27})(t),t)) = (b'_{25})^{(4)}$ then $T_{25} \to \infty$.<br>**Definition of** $(m)^{(4)}$ and $\varepsilon_4$ :<br>Indeed let $t_4$ be so that for $t > t_4$<br>$(b_{25})^{(4)} - (b_i'')^{(4)}((G_{27})(t),t) < \varepsilon_4, T_{24}(t) > (m)^{(4)}$ | 231 |
| Then $\frac{dT_{25}}{dt} \geq (a_{25})^{(4)}(m)^{(4)} - \varepsilon_4 T_{25}$ which leads to<br>$T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{\varepsilon_4}\right)(1 - e^{-\varepsilon_4 t}) + T_{25}^0 e^{-\varepsilon_4 t}$ If we take t such that $e^{-\varepsilon_4 t} = \frac{1}{2}$ it results<br>$T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{2}\right)$, $t = log\frac{2}{\varepsilon_4}$ By taking now $\varepsilon_4$ sufficiently small one sees that $T_{25}$ is unbounded.<br>The same property holds for $T_{26}$ if $\lim_{t\to\infty}(b_{26}'')^{(4)}\left((G_{27})(t),t\right) = (b'_{26})^{(4)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 42<br>Analogous inequalities hold also for $G_{29}$ , $G_{30}, T_{28}, T_{29}, T_{30}$ | 232 |
| It is now sufficient to take $\frac{(a_i)^{(5)}}{(\widehat{M}_{28})^{(5)}}$ , $\frac{(b_i)^{(5)}}{(\widehat{M}_{28})^{(5)}} < 1$ and to choose<br>$(\widehat{P}_{28})^{(5)}$ and $(\widehat{Q}_{28})^{(5)}$ large to have | 233 |
| $\frac{(a_i)^{(5)}}{(\widehat{M}_{28})^{(5)}}\left[(\widehat{P}_{28})^{(5)} + \left((\widehat{P}_{28})^{(5)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{28})^{(5)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{28})^{(5)}$ | 234 |
| $\frac{(b_i)^{(5)}}{(\widehat{M}_{28})^{(5)}}\left[\left((\widehat{Q}_{28})^{(5)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{28})^{(5)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{28})^{(5)}\right] \leq (\widehat{Q}_{28})^{(5)}$ | 235 |
| In order that the operator $\mathcal{A}^{(5)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying | |

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| Equations into itself | |
| The operator $\mathcal{A}^{(5)}$ is a contraction with respect to the metric<br>$d\left(((G_{31})^{(1)},(T_{31})^{(1)}),((G_{31})^{(2)},(T_{31})^{(2)})\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{28})^{(5)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{28})^{(5)}t}\}$<br>Indeed if we denote<br>**Definition of** $(\widetilde{G_{31}}),(\widetilde{T_{31}})$ : $\left((\widetilde{G_{31}}),(\widetilde{T_{31}})\right)=\mathcal{A}^{(5)}((G_{31}),(T_{31}))$<br>It results<br>$\left\|\tilde{G}_{28}^{(1)}-\tilde{G}_i^{(2)}\right\|\le\int_0^t(a_{28})^{(5)}\left\|G_{29}^{(1)}-G_{29}^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{(\widehat{M}_{28})^{(5)}s_{(28)}}\,ds_{(28)}+$<br>$\int_0^t\{(a_{28}^{'})^{(5)}\left\|G_{28}^{(1)}-G_{28}^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}+$<br>$(a_{28}^{''})^{(5)}\left(T_{29}^{(1)},s_{(28)}\right)\left\|G_{28}^{(1)}-G_{28}^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{(\widehat{M}_{28})^{(5)}s_{(28)}}+$<br>$G_{28}^{(2)}\left\|(a_{28}^{''})^{(5)}\left(T_{29}^{(1)},s_{(28)}\right)-(a_{28}^{''})^{(5)}\left(T_{29}^{(2)},s_{(28)}\right)\right\|\ e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{(\widehat{M}_{28})^{(5)}s_{(28)}}\}ds_{(28)}$<br>Where $s_{(28)}$ represents integrand that is integrated over the interval [0, t]<br>From the hypotheses on it follows | 236 |
| $\left\|(G_{31})^{(1)}-(G_{31})^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}t}\le$<br>$\frac{1}{(\widehat{M}_{28})^{(5)}}\left((a_{28})^{(5)}+(a_{28}^{'})^{(5)}+(\widehat{A}_{28})^{(5)}+(\widehat{P}_{28})^{(5)}(\widehat{k}_{28})^{(5)}\right)d\left(((G_{31})^{(1)},(T_{31})^{(1)};\ (G_{31})^{(2)},(T_{31})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 237 |
| **Remark 21:** The fact that we supposed $(a_{28}^{''})^{(5)}$ and $(b_{28}^{''})^{(5)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{28})^{(5)}e^{(\widehat{M}_{28})^{(5)}t}$ *and* $(\widehat{Q}_{28})^{(5)}e^{(\widehat{M}_{28})^{(5)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(5)}$ and $(b_i^{''})^{(5)}, i=28,29,30$ depend only on $T_{29}$ and respectively on $(G_{31})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 238 |
| **Remark 22:** There does not exist any $t$ where $G_i(t)=0$ *and* $T_i(t)=0$<br>it results<br>$G_i(t)\ge G_i^0e^{\left[-\int_0^t\{(a_i^{'})^{(5)}-(a_i^{''})^{(5)}(T_{29}(s_{(28)}),s_{(28)})\}ds_{(28)}\right]}\ge 0$ | 239 |

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| $T_i\,(t) \geq T_i^0 e^{(-(b_i^{'})^{(5)}t)} > 0$ for t > 0 | |
| **Definition of** $\left((\widehat{M}_{28})^{(5)}\right)_1, \left((\widehat{M}_{28})^{(5)}\right)_2$ *and* $\left((\widehat{M}_{28})^{(5)}\right)_3$ :<br><br>**Remark 23:** if $G_{28}$ is bounded, the same property have also $G_{29}$ *and* $G_{30}$ . indeed if<br><br>$G_{28} < (\widehat{M}_{28})^{(5)}$ it follows $\frac{dG_{29}}{dt} \leq \left((\widehat{M}_{28})^{(5)}\right)_1 - (a_{29}^{'})^{(5)} G_{29}$ and by integrating<br><br>$G_{29} \leq \left((\widehat{M}_{28})^{(5)}\right)_2 = G_{29}^0 + 2(a_{29})^{(5)}\left((\widehat{M}_{28})^{(5)}\right)_1/(a_{29}^{'})^{(5)}$<br><br>In the same way , one can obtain<br><br>$G_{30} \leq \left((\widehat{M}_{28})^{(5)}\right)_3 = G_{30}^0 + 2(a_{30})^{(5)}\left((\widehat{M}_{28})^{(5)}\right)_2/(a_{30}^{'})^{(5)}$<br><br>If $G_{29}$ *or* $G_{30}$ is bounded, the same property follows for $G_{28}$ , $G_{30}$ and $G_{28}$ , $G_{29}$ respectively. | 240 |
| **Remark 24:** If $G_{28}$ *is* bounded, from below, the same property holds for $G_{29}$ *and* $G_{30}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{29}$ is bounded from below. | 241 |
| **Remark 25:** If $T_{28}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(5)}((G_{31})(t),t)) = (b_{29}^{'})^{(5)}$ then $T_{29} \to \infty$.<br><br>**Definition of** $(m)^{(5)}$ and $\varepsilon_5$ :<br><br>Indeed let $t_5$ be so that for $t > t_5$<br><br>$(b_{29})^{(5)} - (b_i^{''})^{(5)}((G_{31})(t),t) < \varepsilon_5, T_{28}\,(t) > (m)^{(5)}$ | 242 |
| Then $\frac{dT_{29}}{dt} \geq (a_{29})^{(5)}(m)^{(5)} - \varepsilon_5 T_{29}$ which leads to<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{\varepsilon_5}\right)(1 - e^{-\varepsilon_5 t}) + T_{29}^0 e^{-\varepsilon_5 t}$ If we take t such that $e^{-\varepsilon_5 t} = \frac{1}{2}$ it results<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{2}\right), \quad t = log\frac{2}{\varepsilon_5}$ By taking now $\varepsilon_5$ sufficiently small one sees that $T_{29}$ is unbounded.<br>The same property holds for $T_{30}$ if $\lim_{t\to\infty}(b_{30}^{''})^{(5)}\left((G_{31})(t),t\right) = (b_{30}^{'})^{(5)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations<br><br>Analogous inequalities hold also for $G_{33}$ , $G_{34}, T_{32}, T_{33}, T_{34}$ | 243 |
| It is now sufficient to take $\frac{(a_i)^{(6)}}{(\widehat{M}_{32})^{(6)}}$ , $\frac{(b_i)^{(6)}}{(\widehat{M}_{32})^{(6)}} < 1$ and to choose<br><br>$(\widehat{P}_{32})^{(6)}$ and $(\widehat{Q}_{32})^{(6)}$ large to have | 244 |
| $\frac{(a_i)^{(6)}}{(\widehat{M}_{32})^{(6)}}\left[(\widehat{P}_{32})^{(6)} + \left((\widehat{P}_{32})^{(6)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{32})^{(6)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{32})^{(6)}$ | 245 |

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| $\frac{(b_i)^{(6)}}{(\widehat{M}_{32})^{(6)}}\left[\left((\hat{Q}_{32})^{(6)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{32})^{(6)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{32})^{(6)}\right]\leq(\hat{Q}_{32})^{(6)}$ | 246 |
| In order that the operator $\mathcal{A}^{(6)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(6)}$ is a contraction with respect to the metric<br>$d\left(((G_{35})^{(1)},(T_{35})^{(1)}),((G_{35})^{(2)},(T_{35})^{(2)})\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{32})^{(6)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{32})^{(6)}t}\}$<br>Indeed if we denote<br>**Definition of** $(\widetilde{G_{35}}),(\widetilde{T_{35}})$ : $\left((\widetilde{G_{35}}),(\widetilde{T_{35}})\right)=\mathcal{A}^{(6)}\left((G_{35}),(T_{35})\right)$<br>It results<br>$\left\|\tilde{G}_{32}^{(1)}-\tilde{G}_i^{(2)}\right\|\leq\int_0^t(a_{32})^{(6)}\left\|G_{33}^{(1)}-G_{33}^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{(\widehat{M}_{32})^{(6)}s_{(32)}}ds_{(32)}+$<br>$\int_0^t\{(a'_{32})^{(6)}\left\|G_{32}^{(1)}-G_{32}^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}+$<br>$(a''_{32})^{(6)}\left(T_{33}^{(1)},s_{(32)}\right)\left\|G_{32}^{(1)}-G_{32}^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{(\widehat{M}_{32})^{(6)}s_{(32)}}+$<br>$G_{32}^{(2)}\left\|(a''_{32})^{(6)}\left(T_{33}^{(1)},s_{(32)}\right)-(a''_{32})^{(6)}\left(T_{33}^{(2)},s_{(32)}\right)\right\|\ e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{(\widehat{M}_{32})^{(6)}s_{(32)}}\}ds_{(32)}$<br>Where $s_{(32)}$ represents integrand that is integrated over the interval $[0,t]$<br>From the hypotheses it follows | 247 |
| $\left\|(G_{35})^{(1)}-(G_{35})^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}t}\leq$<br>$\frac{1}{(\widehat{M}_{32})^{(6)}}\left((a_{32})^{(6)}+(a'_{32})^{(6)}+(\widehat{A}_{32})^{(6)}+(\widehat{P}_{32})^{(6)}(\widehat{k}_{32})^{(6)}\right)d\left(((G_{35})^{(1)},(T_{35})^{(1)};\ (G_{35})^{(2)},(T_{35})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 248 |
| **Remark 26:** The fact that we supposed $(a''_{32})^{(6)}$ and $(b''_{32})^{(6)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{32})^{(6)}e^{(\widehat{M}_{32})^{(6)}t}$ *and* $(\widehat{Q}_{32})^{(6)}e^{(\widehat{M}_{32})^{(6)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i'')^{(6)}$ and $(b_i'')^{(6)}, i=32,33,34$ depend only on $T_{33}$ and respectively on $(G_{35})$(*and not on* $t$) and hypothesis can replaced by a usual Lipschitz condition. | 249 |

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| **Remark 27:** There does not exist any $t$ where $G_i(t) = 0$ *and* $T_i(t) = 0$<br><br>it results<br><br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i')^{(6)} - (a_i'')^{(6)}(T_{33}(s_{(32)}), s_{(32)})\} ds_{(32)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{(-(b_i')^{(6)} t)} > 0$ for t > 0 | 250 |
| **Definition of** $\left((\widehat{M}_{32})^{(6)}\right)_1, \left((\widehat{M}_{32})^{(6)}\right)_2$ *and* $\left((\widehat{M}_{32})^{(6)}\right)_3$ :<br><br>**Remark 28:** if $G_{32}$ is bounded, the same property have also $G_{33}$ *and* $G_{34}$ . indeed if<br><br>$G_{32} < (\widehat{M}_{32})^{(6)}$ it follows $\frac{dG_{33}}{dt} \leq \left((\widehat{M}_{32})^{(6)}\right)_1 - (a_{33}')^{(6)} G_{33}$ and by integrating<br><br>$G_{33} \leq \left((\widehat{M}_{32})^{(6)}\right)_2 = G_{33}^0 + 2(a_{33})^{(6)} \left((\widehat{M}_{32})^{(6)}\right)_1 / (a_{33}')^{(6)}$<br><br>In the same way , one can obtain<br><br>$G_{34} \leq \left((\widehat{M}_{32})^{(6)}\right)_3 = G_{34}^0 + 2(a_{34})^{(6)} \left((\widehat{M}_{32})^{(6)}\right)_2 / (a_{34}')^{(6)}$<br><br>If $G_{33}$ *or* $G_{34}$ is bounded, the same property follows for $G_{32}$ , $G_{34}$ and $G_{32}$ , $G_{33}$ respectively. | 251 |
| **Remark 29:** If $G_{32}$ *is* bounded, from below, the same property holds for $G_{33}$ *and* $G_{34}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{33}$ is bounded from below. | 252 |
| **Remark 30:** If $T_{32}$ is bounded from below and $\lim_{t\to\infty} ((b_i'')^{(6)}((G_{35})(t), t)) = (b_{33}')^{(6)}$ then $T_{33} \to \infty$.<br><br>**Definition of** $(m)^{(6)}$ and $\varepsilon_6$ :<br><br>Indeed let $t_6$ be so that for $t > t_6$<br><br>$(b_{33})^{(6)} - (b_i'')^{(6)}((G_{35})(t), t) < \varepsilon_6, T_{32}(t) > (m)^{(6)}$ | 253 |
| Then $\frac{dT_{33}}{dt} \geq (a_{33})^{(6)}(m)^{(6)} - \varepsilon_6 T_{33}$ which leads to<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{\varepsilon_6}\right)(1 - e^{-\varepsilon_6 t}) + T_{33}^0 e^{-\varepsilon_6 t}$ If we take t such that $e^{-\varepsilon_6 t} = \frac{1}{2}$ it results<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{2}\right)$, $t = log\frac{2}{\varepsilon_6}$ By taking now $\varepsilon_6$ sufficiently small one sees that $T_{33}$ is unbounded.<br>The same property holds for $T_{34}$ if $\lim_{t\to\infty} (b_{34}'')^{(6)}\left((G_{35})(t), t(t), t\right) = (b_{34}')^{(6)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 254 |
| Analogous inequalities hold also for $G_{37}$ , $G_{38}, T_{36}, T_{37}, T_{38}$<br><br>It is now sufficient to take $\frac{(a_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}$ , $\frac{(b_i)^{(7)}}{(\widehat{M}_{36})^{(7)}} < 1$ and to choose<br>$(\widehat{P}_{36})^{(7)}$ and $(\widehat{Q}_{36})^{(7)}$ large to have | 255 |

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| $\frac{(a_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}\left[(\widehat{P}_{36})^{(7)}+\left((\hat{P}_{36})^{(7)}+G_j^0\right)e^{-\left(\frac{(\hat{P}_{36})^{(7)}+G_j^0}{G_j^0}\right)}\right]\leq(\hat{P}_{36})^{(7)}$ | 256 |
| $\frac{(b_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}\left[\left((\hat{Q}_{36})^{(7)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{36})^{(7)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{36})^{(7)}\right]\leq(\hat{Q}_{36})^{(7)}$ | 257 |
| In order that the operator $\mathcal{A}^{(7)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(7)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{39})^{(1)},(T_{39})^{(1)}),((G_{39})^{(2)},(T_{39})^{(2)})\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{36})^{(7)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{36})^{(7)}t}\}$<br><br>Indeed if we denote<br><br>**Definition of** $\widetilde{(G_{39})},\widetilde{(T_{39})}:\left(\widetilde{(G_{39})},\widetilde{(T_{39})}\right)=\mathcal{A}^{(7)}((G_{39}),(T_{39}))$<br>It results<br><br>$\left\|\tilde{G}_{36}^{(1)}-\tilde{G}_i^{(2)}\right\|\leq\int_0^t(a_{36})^{(7)}\left\|G_{37}^{(1)}-G_{37}^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{(\widehat{M}_{36})^{(7)}s_{(36)}}\,ds_{(36)}+$<br>$\int_0^t\{(a_{36}')^{(7)}\left\|G_{36}^{(1)}-G_{36}^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}+$<br>$(a_{36}'')^{(7)}\left(T_{37}^{(1)},s_{(36)}\right)\left\|G_{36}^{(1)}-G_{36}^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{(\widehat{M}_{36})^{(7)}s_{(36)}}+$<br>$G_{36}^{(2)}\left\|(a_{36}'')^{(7)}\left(T_{37}^{(1)},s_{(36)}\right)-(a_{36}'')^{(7)}\left(T_{37}^{(2)},s_{(36)}\right)\right\|\ e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{(\widehat{M}_{36})^{(7)}s_{(36)}}\}ds_{(36)}$<br><br>Where $s_{(36)}$ represents integrand that is integrated over the interval $[0,t]$<br><br>From the hypotheses on it follows | 258 |
| $\left\|(G_{39})^{(1)}-(G_{39})^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}t}\leq$<br>$\frac{1}{(\widehat{M}_{36})^{(7)}}\left((a_{36})^{(7)}+(a_{36}')^{(7)}+(\widehat{A}_{36})^{(7)}+(\widehat{P}_{36})^{(7)}(\widehat{k}_{36})^{(7)}\right)d\left(((G_{39})^{(1)},(T_{39})^{(1)};\ (G_{39})^{(2)},(T_{39})^{(2)})\right)$<br><br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 259 |
| **Remark 31:** The fact that we supposed $(a_{36}'')^{(7)}$ and $(b_{36}'')^{(7)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{36})^{(7)}e^{(\widehat{M}_{36})^{(7)}t}$ *and* $(\widehat{Q}_{36})^{(7)}e^{(\widehat{M}_{36})^{(7)}t}$ respectively of $\mathbb{R}_+$. | 260 |

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| If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(7)}$ and $(b_i^{''})^{(7)}, i = 36,37,38$ depend only on $T_{37}$ and respectively on $(G_{39})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 32:** There does not exist any $t$ where $G_i\,(t) = 0\ and\ T_i\,(t) = 0$<br><br>it results<br><br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(7)} - (a_i^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\} ds_{(36)}\right]} \geq 0$<br><br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(7)} t\right)} > 0$ for t > 0 | 261 |
| **Definition of** $\left((\widehat{M}_{36})^{(7)}\right)_1, \left((\widehat{M}_{36})^{(7)}\right)_2\ and\ \left((\widehat{M}_{36})^{(7)}\right)_3$ :<br><br>**Remark 33:** if $G_{36}$ is bounded, the same property have also $G_{37}\ and\ G_{38}$ . indeed if<br><br>$G_{36} < (\widehat{M}_{36})^{(7)}$ it follows $\frac{dG_{37}}{dt} \leq \left((\widehat{M}_{36})^{(7)}\right)_1 - (a_{37}^{'})^{(7)} G_{37}$ and by integrating<br><br>$G_{37} \leq \left((\widehat{M}_{36})^{(7)}\right)_2 = G_{37}^0 + 2(a_{37})^{(7)}\left((\widehat{M}_{36})^{(7)}\right)_1 / (a_{37}^{'})^{(7)}$<br><br>In the same way , one can obtain<br><br>$G_{38} \leq \left((\widehat{M}_{36})^{(7)}\right)_3 = G_{38}^0 + 2(a_{38})^{(7)}\left((\widehat{M}_{36})^{(7)}\right)_2 / (a_{38}^{'})^{(7)}$<br><br>If $G_{37}\ or\ G_{38}$ is bounded, the same property follows for $G_{36}$ , $G_{38}$ and $G_{36}$ , $G_{37}$ respectively. | 262 |
| **Remark 34:** If $G_{36}$ *is* bounded, from below, the same property holds for $G_{37}\ and\ G_{38}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{37}$ is bounded from below. | 263 |
| **Remark 35:** If $T_{36}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(7)}((G_{39})(t), t)) = (b_{37}^{'})^{(7)}$ then $T_{37} \to \infty$.<br><br>**Definition of** $(m)^{(7)}$ and $\varepsilon_7$ :<br><br>Indeed let $t_7$ be so that for $t > t_7$<br><br>$(b_{37})^{(7)} - (b_i^{''})^{(7)}((G_{39})(t), t) < \varepsilon_7, T_{36}\,(t) > (m)^{(7)}$ | 264 |
| Then $\frac{dT_{37}}{dt} \geq (a_{37})^{(7)}(m)^{(7)} - \varepsilon_7 T_{37}$ which leads to<br><br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{\varepsilon_7}\right)(1 - e^{-\varepsilon_7 t}) + T_{37}^0 e^{-\varepsilon_7 t}$ If we take t such that $e^{-\varepsilon_7 t} = \frac{1}{2}$ it results<br><br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{2}\right),\ \ t = log\frac{2}{\varepsilon_7}$ By taking now $\varepsilon_7$ sufficiently small one sees that $T_{37}$ is unbounded. | 265 |

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| The same property holds for $T_{38}$ if $\lim_{t\to\infty}(b_{38}^{''})^{(7)}\left((G_{39})(t),t\right)=(b_{38}^{'})^{(7)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |
| It is now sufficient to take $\frac{(a_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}$ , $\frac{(b_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}<1$ and to choose<br>$(\widehat{P}_{40})^{(8)}$ and $(\widehat{Q}_{40})^{(8)}$ large to have | 266 |
| $\frac{(a_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}\left[(\widehat{P}_{40})^{(8)}+\left((\hat{P}_{40})^{(8)}+G_j^0\right)e^{-\left(\frac{(\hat{P}_{40})^{(8)}+G_j^0}{G_j^0}\right)}\right]\leq(\hat{P}_{40})^{(8)}$ | 267 |
| $\frac{(b_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}\left[\left((\hat{Q}_{40})^{(8)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{40})^{(8)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{40})^{(8)}\right]\leq(\hat{Q}_{40})^{(8)}$ | 268 |
| In order that the operator $\mathcal{A}^{(8)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(8)}$ is a contraction with respect to the metric | |
| $d\left(((G_{43})^{(1)},(T_{43})^{(1)}),((G_{43})^{(2)},(T_{43})^{(2)})\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{40})^{(8)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{40})^{(8)}t}\}$ | 269 |
| Indeed if we denote<br><br>**Definition of** $\widetilde{(G_{43})},\widetilde{(T_{43})}$ : $\left(\widetilde{(G_{43})},\widetilde{(T_{43})}\right)=\mathcal{A}^{(8)}((G_{43}),(T_{43}))$ | 270 |
| It results<br><br>$\left\|\tilde{G}_{40}^{(1)}-\tilde{G}_i^{(2)}\right\|\leq\int_0^t(a_{40})^{(8)}\left\|G_{41}^{(1)}-G_{41}^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{(\widehat{M}_{40})^{(8)}s_{(40)}}ds_{(40)}+$<br>$\int_0^t\{(a_{40}^{'})^{(8)}\left\|G_{40}^{(1)}-G_{40}^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}+$<br>$(a_{40}^{''})^{(8)}\left(T_{41}^{(1)},s_{(40)}\right)\left\|G_{40}^{(1)}-G_{40}^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{(\widehat{M}_{40})^{(8)}s_{(40)}}+$<br>$G_{40}^{(2)}\left\\|(a_{40}^{''})^{(8)}\left(T_{41}^{(1)},s_{(40)}\right)-(a_{40}^{''})^{(8)}\left(T_{41}^{(2)},s_{(40)}\right)\right\\|\ e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{(\widehat{M}_{40})^{(8)}s_{(40)}}\}ds_{(40)}$ | 271 |
| Where $s_{(40)}$ represents integrand that is integrated over the interval $[0,t]$<br><br>From the hypotheses it follows | 272 |
| $\left\|(G_{43})^{(1)}-(G_{43})^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}t}\leq$<br>$\frac{1}{(\widehat{M}_{40})^{(8)}}\left((a_{40})^{(8)}+(a_{40}^{'})^{(8)}+(\widehat{A}_{40})^{(8)}+(\widehat{P}_{40})^{(8)}(\widehat{k}_{40})^{(8)}\right)d\left(((G_{43})^{(1)},(T_{43})^{(1)};\ (G_{43})^{(2)},(T_{43})^{(2)})\right)$<br><br>And analogous inequalities for $G_i$ $and$ $T_i$. Taking into account the hypothesis the result follows | 273 |

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| **Remark 36:** The fact that we supposed $(a_{40}^{''})^{(8)}$ and $(b_{40}^{''})^{(8)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{40})^{(8)}e^{(\widehat{M}_{40})^{(8)}t}$ $and$ $(\widehat{Q}_{40})^{(8)}e^{(\widehat{M}_{40})^{(8)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(8)}$ and $(b_i^{''})^{(8)}, i = 40,41,42$ depend only on $T_{41}$ and respectively on $(G_{43})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 274 |
| **Remark 37** There does not exist any $t$ where $G_i\,(t) = 0\ and\ T_i\,(t) = 0$<br><br>it results<br><br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(8)} - (a_i^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\} ds_{(40)}\right]} \geq 0$<br><br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(8)}t\right)} > 0$ for t > 0 | 275 |
| **Definition of** $\left((\widehat{M}_{40})^{(8)}\right)_1, \left((\widehat{M}_{40})^{(8)}\right)_2\ and\ \left((\widehat{M}_{40})^{(8)}\right)_3$ :<br><br>**Remark 38:** if $G_{40}$ is bounded, the same property have also $G_{41}\ and\ G_{42}$ . indeed if<br><br>$G_{40} < (\widehat{M}_{40})^{(8)}$ it follows $\frac{dG_{41}}{dt} \leq \left((\widehat{M}_{40})^{(8)}\right)_1 - (a_{41}^{'})^{(8)} G_{41}$ and by integrating<br><br>$G_{41} \leq \left((\widehat{M}_{40})^{(8)}\right)_2 = G_{41}^0 + 2(a_{41})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_1 / (a_{41}^{'})^{(8)}$<br><br>In the same way , one can obtain<br><br>$G_{42} \leq \left((\widehat{M}_{40})^{(8)}\right)_3 = G_{42}^0 + 2(a_{42})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_2 / (a_{42}^{'})^{(8)}$<br><br>If $G_{41}\ or\ G_{42}$ is bounded, the same property follows for $G_{40}$ , $G_{42}$ and $G_{40}$ , $G_{41}$ respectively. | 276 |
| **Remark 39:** If $G_{40}$ $is$ bounded, from below, the same property holds for $G_{41}\ and\ G_{42}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{41}$ is bounded from below. | 277 |
| **Remark 40:** If $T_{40}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(8)}((G_{43})(t), t)) = (b_{41}^{'})^{(8)}$ then $T_{41} \to \infty$.<br><br>**Definition of** $(m)^{(8)}$ and $\varepsilon_8$ :<br><br>Indeed let $t_8$ be so that for $t > t_8$<br><br>$(b_{41})^{(8)} - (b_i^{''})^{(8)}\left((G_{43})(t), t\right) < \varepsilon_8, T_{40}\,(t) > (m)^{(8)}$ | 278 |

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| Then $\frac{dT_{41}}{dt} \geq (a_{41})^{(8)}(m)^{(8)} - \varepsilon_8 T_{41}$ which leads to<br><br>$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{\varepsilon_8}\right)(1 - e^{-\varepsilon_8 t}) + T_{41}^0 e^{-\varepsilon_8 t}$ If we take t such that $e^{-\varepsilon_8 t} = \frac{1}{2}$ it results<br><br>$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{2}\right), \quad t = log\frac{2}{\varepsilon_8}$ By taking now $\varepsilon_8$ sufficiently small one sees that $T_{41}$ is unbounded.<br>The same property holds for $T_{42}$ if $\lim_{t\to\infty}(b_{42}^{''})^{(8)}\left((G_{43})(t), t(t), t\right) = (b_{42}^{'})^{(8)}$ | 279 |
| It is now sufficient to take $\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}, \frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}} < 1$ and to choose $(\widehat{P}_{44})^{(9)}$ and $(\widehat{Q}_{44})^{(9)}$ large to have | 279 A |
| $\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[(\widehat{P}_{44})^{(9)} + \left((\widehat{P}_{44})^{(9)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{44})^{(9)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{44})^{(9)}$ | |
| $\frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[\left((\widehat{Q}_{44})^{(9)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{44})^{(9)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{44})^{(9)}\right] \leq (\widehat{Q}_{44})^{(9)}$ | |
| In order that the operator $\mathcal{A}^{(9)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying 39,35,36 into itself | |
| The operator $\mathcal{A}^{(9)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{47})^{(1)}, (T_{47})^{(1)}), ((G_{47})^{(2)}, (T_{47})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t) - G_i^{(2)}(t)\right|e^{-(\widehat{M}_{44})^{(9)}t}, \max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t) - T_i^{(2)}(t)\right|e^{-(\widehat{M}_{44})^{(9)}t}\}$<br><br>Indeed if we denote<br><br>**Definition of** $\widetilde{(G_{47})}, \widetilde{(T_{47})} : \left(\widetilde{(G_{47})}, \widetilde{(T_{47})}\right) = \mathcal{A}^{(9)}\left((G_{47}), (T_{47})\right)$<br>It results<br><br>$\left\|\tilde{G}_{44}^{(1)} - \tilde{G}_i^{(2)}\right| \leq \int_0^t (a_{44})^{(9)}\left|G_{45}^{(1)} - G_{45}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\,ds_{(44)} +$<br>$\int_0^t\{(a_{44}^{'})^{(9)}\left|G_{44}^{(1)} - G_{44}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{-(\widehat{M}_{44})^{(9)}s_{(44)}} +$<br>$(a_{44}^{''})^{(9)}\left(T_{45}^{(1)}, s_{(44)}\right)\left|G_{44}^{(1)} - G_{44}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}} +$<br>$G_{44}^{(2)}\left|(a_{44}^{''})^{(9)}\left(T_{45}^{(1)}, s_{(44)}\right) - (a_{44}^{''})^{(9)}\left(T_{45}^{(2)}, s_{(44)}\right)\right| \; e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\}ds_{(44)}$<br><br>Where $s_{(44)}$ represents integrand that is integrated over the interval $[0, t]$<br><br>From the hypotheses on 45,46,47,28 and 29 it follows | |
| $\left|(G_{47})^{(1)} - G^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}t} \leq$<br>$\frac{1}{(\widehat{M}_{44})^{(9)}}\left((a_{44})^{(9)} + (a_{44}^{'})^{(9)} + (\widehat{A}_{44})^{(9)} + (\widehat{P}_{44})^{(9)}(\widehat{k}_{44})^{(9)}\right)d\left(((G_{47})^{(1)}, (T_{47})^{(1)}; (G_{47})^{(2)}, (T_{47})^{(2)})\right)$ | |

And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis (39,35,36) the result follows

**Remark 41:** The fact that we supposed $(a_{44}^{''})^{(9)}$ and $(b_{44}^{''})^{(9)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{44})^{(9)}e^{(\widehat{M}_{44})^{(9)}t}$ *and* $(\widehat{Q}_{44})^{(9)}e^{(\widehat{M}_{44})^{(9)}t}$ respectively of $\mathbb{R}_+$.

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(9)}$ and $(b_i^{''})^{(9)}, i = 44{,}45{,}46$ depend only on $T_{45}$ and respectively on $(G_{47})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition.

**Remark 42:** There does not exist any $t$ where $G_i\,(t) = 0\ and\ T_i\,(t) = 0$

From 99 to 44 it results

$$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(9)} - (a_i^{''})^{(9)}(T_{45}(s_{(44)}), s_{(44)})\} ds_{(44)}\right]} \geq 0$$

$$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(9)}t\right)} > 0 \quad \text{for t} > 0$$

**Definition of** $\left((\widehat{M}_{44})^{(9)}\right)_1, \left((\widehat{M}_{44})^{(9)}\right)_2\ and\ \left((\widehat{M}_{44})^{(9)}\right)_3$ :

**Remark 43:** if $G_{44}$ is bounded, the same property have also $G_{45}\ and\ G_{46}$ . indeed if

$G_{44} < (\widehat{M}_{44})^{(9)}$ it follows $\frac{dG_{45}}{dt} \leq \left((\widehat{M}_{44})^{(9)}\right)_1 - (a_{45}^{'})^{(9)}G_{45}$ and by integrating

$$G_{45} \leq \left((\widehat{M}_{44})^{(9)}\right)_2 = G_{45}^0 + 2(a_{45})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_1/(a_{45}^{'})^{(9)}$$

In the same way , one can obtain

$$G_{46} \leq \left((\widehat{M}_{44})^{(9)}\right)_3 = G_{46}^0 + 2(a_{46})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_2/(a_{46}^{'})^{(9)}$$

If $G_{45}\ or\ G_{46}$ is bounded, the same property follows for $G_{44}$ , $G_{46}$ and $G_{44}$ , $G_{45}$ respectively.

**Remark 44:** If $G_{44}$ *is* bounded, from below, the same property holds for $G_{45}\ and\ G_{46}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{45}$ is bounded from below.

**Remark 45:** If $T_{44}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(9)}\left((G_{47})(t), t\right) = (b_{45}^{'})^{(9)}$ then $T_{45} \to \infty$.

**Definition of** $(m)^{(9)}$ and $\varepsilon_9$ :

Indeed let $t_9$ be so that for $t > t_9$

$$(b_{45})^{(9)} - (b_i^{''})^{(9)}\left((G_{47})(t), t\right) < \varepsilon_9, T_{44}\,(t) > (m)^{(9)}$$

Then $\frac{dT_{45}}{dt} \geq (a_{45})^{(9)}(m)^{(9)} - \varepsilon_9 T_{45}$ which leads to

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{\varepsilon_9}\right)(1 - e^{-\varepsilon_9 t}) + T_{45}^0 e^{-\varepsilon_9 t}$ If we take t such that $e^{-\varepsilon_9 t} = \frac{1}{2}$ it results

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{2}\right)$, $t = log\frac{2}{\varepsilon_9}$ By taking now $\varepsilon_9$ sufficiently small one sees that $T_{45}$ is unbounded. The same property holds for $T_{46}$ if $\lim_{t\to\infty}(b''_{46})^{(9)}\left((G_{47})(t),t\right) = (b'_{46})^{(9)}$
We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 92

**Behavior of the solutions of equation** 280

**Theorem** If we denote and define

**Definition of** $(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ :

$(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ four constants satisfying
$-(\sigma_2)^{(1)} \leq -(a'_{13})^{(1)} + (a'_{14})^{(1)} - (a''_{13})^{(1)}(T_{14}, t) + (a''_{14})^{(1)}(T_{14}, t) \leq -(\sigma_1)^{(1)}$

$-(\tau_2)^{(1)} \leq -(b'_{13})^{(1)} + (b'_{14})^{(1)} - (b''_{13})^{(1)}(G,t) - (b''_{14})^{(1)}(G,t) \leq -(\tau_1)^{(1)}$

**Definition of** $(\nu_1)^{(1)}, (\nu_2)^{(1)}, (u_1)^{(1)}, (u_2)^{(1)}, \nu^{(1)}, u^{(1)}$ : 281

By $(\nu_1)^{(1)} > 0$, $(\nu_2)^{(1)} < 0$ and respectively $(u_1)^{(1)} > 0$, $(u_2)^{(1)} < 0$ the roots of the equations $(a_{14})^{(1)}\left(\nu^{(1)}\right)^2 + (\sigma_1)^{(1)}\nu^{(1)} - (a_{13})^{(1)} = 0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2 + (\tau_1)^{(1)}u^{(1)} - (b_{13})^{(1)} = 0$

**Definition of** $(\bar{\nu}_1)^{(1)},, (\bar{\nu}_2)^{(1)}, (\bar{u}_1)^{(1)}, (\bar{u}_2)^{(1)}$ : 282

By $(\bar{\nu}_1)^{(1)} > 0$, $(\bar{\nu}_2)^{(1)} < 0$ and respectively $(\bar{u}_1)^{(1)} > 0$, $(\bar{u}_2)^{(1)} < 0$ the roots of the equations $(a_{14})^{(1)}\left(\nu^{(1)}\right)^2 + (\sigma_2)^{(1)}\nu^{(1)} - (a_{13})^{(1)} = 0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2 + (\tau_2)^{(1)}u^{(1)} - (b_{13})^{(1)} = 0$

**Definition of** $(m_1)^{(1)}, (m_2)^{(1)}, (\mu_1)^{(1)}, (\mu_2)^{(1)}, (\nu_0)^{(1)}$ :- 283

If we define $(m_1)^{(1)}, (m_2)^{(1)}, (\mu_1)^{(1)}, (\mu_2)^{(1)}$ by
$(m_2)^{(1)} = (\nu_0)^{(1)}, (m_1)^{(1)} = (\nu_1)^{(1)}$, $if\ (\nu_0)^{(1)} < (\nu_1)^{(1)}$

$(m_2)^{(1)} = (\nu_1)^{(1)}, (m_1)^{(1)} = (\bar{\nu}_1)^{(1)}, if\ (\nu_1)^{(1)} < (\nu_0)^{(1)} < (\bar{\nu}_1)^{(1)}$,

and $\boxed{(\nu_0)^{(1)} = \frac{G^0_{13}}{G^0_{14}}}$

$(m_2)^{(1)} = (\nu_1)^{(1)}, (m_1)^{(1)} = (\nu_0)^{(1)}$, $if\ (\bar{\nu}_1)^{(1)} < (\nu_0)^{(1)}$

and analogously 284

$(\mu_2)^{(1)} = (u_0)^{(1)}, (\mu_1)^{(1)} = (u_1)^{(1)}$, $if\ (u_0)^{(1)} < (u_1)^{(1)}$

$(\mu_2)^{(1)} = (u_1)^{(1)}, (\mu_1)^{(1)} = (\bar{u}_1)^{(1)}, if\ (u_1)^{(1)} < (u_0)^{(1)} < (\bar{u}_1)^{(1)}$,

and $\boxed{(u_0)^{(1)} = \frac{T^0_{13}}{T^0_{14}}}$

$(\mu_2)^{(1)} = (u_1)^{(1)}, (\mu_1)^{(1)} = (u_0)^{(1)}, if\ (\bar{u}_1)^{(1)} < (u_0)^{(1)}$ where $(u_1)^{(1)}, (\bar{u}_1)^{(1)}$

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| are defined | |
| Then the solution of global equations satisfies the inequalities<br>$G_{13}^0 e^{((S_1)^{(1)}-(p_{13})^{(1)})t} \le G_{13}(t) \le G_{13}^0 e^{(S_1)^{(1)}t}$<br>where $(p_i)^{(1)}$ is defined by equation<br>$\frac{1}{(m_1)^{(1)}} G_{13}^0 e^{((S_1)^{(1)}-(p_{13})^{(1)})t} \le G_{14}(t) \le \frac{1}{(m_2)^{(1)}} G_{13}^0 e^{(S_1)^{(1)}t}$ | 285 |
| $\left(\frac{(a_{15})^{(1)}G_{13}^0}{(m_1)^{(1)}((S_1)^{(1)}-(p_{13})^{(1)}-(S_2)^{(1)})}\left[e^{((S_1)^{(1)}-(p_{13})^{(1)})t} - e^{-(S_2)^{(1)}t}\right] + G_{15}^0 e^{-(S_2)^{(1)}t} \le G_{15}(t) \le \right.$<br>$\left.\frac{(a_{15})^{(1)}G_{13}^0}{(m_2)^{(1)}((S_1)^{(1)}-(a'_{15})^{(1)})}[e^{(S_1)^{(1)}t} - e^{-(a'_{15})^{(1)}t}] + G_{15}^0 e^{-(a'_{15})^{(1)}t}\right)$ | 286 |
| $\boxed{T_{13}^0 e^{(R_1)^{(1)}t} \le T_{13}(t) \le T_{13}^0 e^{((R_1)^{(1)}+(r_{13})^{(1)})t}}$ | 287 |
| $\frac{1}{(\mu_1)^{(1)}} T_{13}^0 e^{(R_1)^{(1)}t} \le T_{13}(t) \le \frac{1}{(\mu_2)^{(1)}} T_{13}^0 e^{((R_1)^{(1)}+(r_{13})^{(1)})t}$ | 288 |
| $\frac{(b_{15})^{(1)}T_{13}^0}{(\mu_1)^{(1)}((R_1)^{(1)}-(b'_{15})^{(1)})}\left[e^{(R_1)^{(1)}t} - e^{-(b'_{15})^{(1)}t}\right] + T_{15}^0 e^{-(b'_{15})^{(1)}t} \le T_{15}(t) \le$<br>$\frac{(a_{15})^{(1)}T_{13}^0}{(\mu_2)^{(1)}((R_1)^{(1)}+(r_{13})^{(1)}+(R_2)^{(1)})}\left[e^{((R_1)^{(1)}+(r_{13})^{(1)})t} - e^{-(R_2)^{(1)}t}\right] + T_{15}^0 e^{-(R_2)^{(1)}t}$ | 289 |
| **Definition of** $(S_1)^{(1)}, (S_2)^{(1)}, (R_1)^{(1)}, (R_2)^{(1)}$:-<br>Where $(S_1)^{(1)} = (a_{13})^{(1)}(m_2)^{(1)} - (a'_{13})^{(1)}$<br>$(S_2)^{(1)} = (a_{15})^{(1)} - (p_{15})^{(1)}$<br>$(R_1)^{(1)} = (b_{13})^{(1)}(\mu_2)^{(1)} - (b'_{13})^{(1)}$<br>$(R_2)^{(1)} = (b'_{15})^{(1)} - (r_{15})^{(1)}$ | 290 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define | 291 |
| **Definition of** $(\sigma_1)^{(2)}, (\sigma_2)^{(2)}, (\tau_1)^{(2)}, (\tau_2)^{(2)}$ :<br>$(\sigma_1)^{(2)}, (\sigma_2)^{(2)}, (\tau_1)^{(2)}, (\tau_2)^{(2)}$ four constants satisfying | 292 |
| $-(\sigma_2)^{(2)} \le -(a'_{16})^{(2)} + (a'_{17})^{(2)} - (a''_{16})^{(2)}(T_{17}, t) + (a''_{17})^{(2)}(T_{17}, t) \le -(\sigma_1)^{(2)}$ | 293 |
| $-(\tau_2)^{(2)} \le -(b'_{16})^{(2)} + (b'_{17})^{(2)} - (b''_{16})^{(2)}((G_{19}), t) - (b''_{17})^{(2)}((G_{19}), t) \le -(\tau_1)^{(2)}$ | 294 |
| **Definition of** $(v_1)^{(2)}, (v_2)^{(2)}, (u_1)^{(2)}, (u_2)^{(2)}$ : | 295 |

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| By $(\nu_1)^{(2)} > 0$ , $(\nu_2)^{(2)} < 0$ and respectively $(u_1)^{(2)} > 0$ , $(u_2)^{(2)} < 0$ the roots | 296 |
| of the equations $(a_{17})^{(2)}\left(\nu^{(2)}\right)^2 + (\sigma_1)^{(2)}\nu^{(2)} - (a_{16})^{(2)} = 0$ | 297 |
| and $(b_{14})^{(2)}\left(u^{(2)}\right)^2 + (\tau_1)^{(2)}u^{(2)} - (b_{16})^{(2)} = 0$ and | 298 |
| **Definition of** $(\bar{\nu}_1)^{(2)}$,, $(\bar{\nu}_2)^{(2)}, (\bar{u}_1)^{(2)}, (\bar{u}_2)^{(2)}$ : | 299 |
| By $(\bar{\nu}_1)^{(2)} > 0$ , $(\bar{\nu}_2)^{(2)} < 0$ and respectively $(\bar{u}_1)^{(2)} > 0$ , $(\bar{u}_2)^{(2)} < 0$ the | 300 |
| roots of the equations $(a_{17})^{(2)}\left(\nu^{(2)}\right)^2 + (\sigma_2)^{(2)}\nu^{(2)} - (a_{16})^{(2)} = 0$ | 301 |
| and $(b_{17})^{(2)}\left(u^{(2)}\right)^2 + (\tau_2)^{(2)}u^{(2)} - (b_{16})^{(2)} = 0$ | 302 |
| **Definition of** $(m_1)^{(2)}$ , $(m_2)^{(2)}$ , $(\mu_1)^{(2)}, (\mu_2)^{(2)}$ :- | 303 |
| If we define $(m_1)^{(2)}$ , $(m_2)^{(2)}$ , $(\mu_1)^{(2)}, (\mu_2)^{(2)}$ by | 304 |
| $(m_2)^{(2)} = (\nu_0)^{(2)}, (m_1)^{(2)} = (\nu_1)^{(2)}$, ***if*** $(\nu_0)^{(2)} < (\nu_1)^{(2)}$ | 305 |
| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ , ***if*** $(\nu_1)^{(2)} < (\nu_0)^{(2)} < (\bar{\nu}_1)^{(2)}$, and $\boxed{(\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0}}$ | 306 |
| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\nu_0)^{(2)}$, ***if*** $(\bar{\nu}_1)^{(2)} < (\nu_0)^{(2)}$ | 307 |
| and analogously $(\mu_2)^{(2)} = (u_0)^{(2)}, (\mu_1)^{(2)} = (u_1)^{(2)}$, ***if*** $(u_0)^{(2)} < (u_1)^{(2)}$ $(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (\bar{u}_1)^{(2)}$ , ***if*** $(u_1)^{(2)} < (u_0)^{(2)} < (\bar{u}_1)^{(2)}$, and $\boxed{(u_0)^{(2)} = \frac{T_{16}^0}{T_{17}^0}}$ | 308 |
| $(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (u_0)^{(2)}$, ***if*** $(\bar{u}_1)^{(2)} < (u_0)^{(2)}$ | 309 |
| Then the solution of global equations satisfies the inequalities $G_{16}^0 e^{((S_1)^{(2)} - (p_{16})^{(2)})t} \le G_{16}(t) \le G_{16}^0 e^{(S_1)^{(2)}t}$ | 310 |
| $(p_i)^{(2)}$ is defined by equation | |
| $\frac{1}{(m_1)^{(2)}} G_{16}^0 e^{((S_1)^{(2)} - (p_{16})^{(2)})t} \le G_{17}(t) \le \frac{1}{(m_2)^{(2)}} G_{16}^0 e^{(S_1)^{(2)}t}$ | 311 |
| $\left(\frac{(a_{18})^{(2)} G_{16}^0}{(m_1)^{(2)}((S_1)^{(2)} - (p_{16})^{(2)} - (S_2)^{(2)})}\left[e^{((S_1)^{(2)} - (p_{16})^{(2)})t} - e^{-(S_2)^{(2)}t}\right] + G_{18}^0 e^{-(S_2)^{(2)}t} \le G_{18}(t) \le\right.$ | 312 |

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| $\frac{(a_{18})^{(2)}G_{16}^{0}}{(m_2)^{(2)}((S_1)^{(2)}-(a_{18}')^{(2)})}[e^{(S_1)^{(2)}t}-e^{-(a_{18}')^{(2)}t}]+G_{18}^{0}e^{-(a_{18}')^{(2)}t})$ | |
| $\boxed{T_{16}^{0}e^{(R_1)^{(2)}t}\le T_{16}(t)\le T_{16}^{0}e^{((R_1)^{(2)}+(r_{16})^{(2)})t}}$ | 313 |
| $\frac{1}{(\mu_1)^{(2)}}T_{16}^{0}e^{(R_1)^{(2)}t}\le T_{16}(t)\le\frac{1}{(\mu_2)^{(2)}}T_{16}^{0}e^{((R_1)^{(2)}+(r_{16})^{(2)})t}$ | 314 |
| $\frac{(b_{18})^{(2)}T_{16}^{0}}{(\mu_1)^{(2)}((R_1)^{(2)}-(b_{18}')^{(2)})}\left[e^{(R_1)^{(2)}t}-e^{-(b_{18}')^{(2)}t}\right]+T_{18}^{0}e^{-(b_{18}')^{(2)}t}\le T_{18}(t)\le$ $\frac{(a_{18})^{(2)}T_{16}^{0}}{(\mu_2)^{(2)}((R_1)^{(2)}+(r_{16})^{(2)}+(R_2)^{(2)})}\left[e^{((R_1)^{(2)}+(r_{16})^{(2)})t}-e^{-(R_2)^{(2)}t}\right]+T_{18}^{0}e^{-(R_2)^{(2)}t}$ | 315 |
| **Definition of** $(S_1)^{(2)},(S_2)^{(2)},(R_1)^{(2)},(R_2)^{(2)}$:- | 316 |
| Where $(S_1)^{(2)}=(a_{16})^{(2)}(m_2)^{(2)}-(a_{16}')^{(2)}$ $(S_2)^{(2)}=(a_{18})^{(2)}-(p_{18})^{(2)}$ | 317 |
| $(R_1)^{(2)}=(b_{16})^{(2)}(\mu_2)^{(1)}-(b_{16}')^{(2)}$ $(R_2)^{(2)}=(b_{18}')^{(2)}-(r_{18})^{(2)}$ | 318 |
| **Behavior of the solutions** **Theorem 3:** If we denote and define **Definition of** $(\sigma_1)^{(3)},(\sigma_2)^{(3)},(\tau_1)^{(3)},(\tau_2)^{(3)}$: $(\sigma_1)^{(3)},(\sigma_2)^{(3)},(\tau_1)^{(3)},(\tau_2)^{(3)}$ four constants satisfying $-(\sigma_2)^{(3)}\le-(a_{20}')^{(3)}+(a_{21}')^{(3)}-(a_{20}'')^{(3)}(T_{21},t)+(a_{21}'')^{(3)}(T_{21},t)\le-(\sigma_1)^{(3)}$ $-(\tau_2)^{(3)}\le-(b_{20}')^{(3)}+(b_{21}')^{(3)}-(b_{20}'')^{(3)}(G_{23},t)-(b_{21}'')^{(3)}((G_{23}),t)\le-(\tau_1)^{(3)}$ | 319 |
| **Definition of** $(v_1)^{(3)},(v_2)^{(3)},(u_1)^{(3)},(u_2)^{(3)}$: By $(v_1)^{(3)}>0,(v_2)^{(3)}<0$ and respectively $(u_1)^{(3)}>0,(u_2)^{(3)}<0$ the roots of the equations $(a_{21})^{(3)}(v^{(3)})^2+(\sigma_1)^{(3)}v^{(3)}-(a_{20})^{(3)}=0$ and $(b_{21})^{(3)}(u^{(3)})^2+(\tau_1)^{(3)}u^{(3)}-(b_{20})^{(3)}=0$ and By $(\bar{v}_1)^{(3)}>0,(\bar{v}_2)^{(3)}<0$ and respectively $(\bar{u}_1)^{(3)}>0,(\bar{u}_2)^{(3)}<0$ the roots of the equations $(a_{21})^{(3)}(v^{(3)})^2+(\sigma_2)^{(3)}v^{(3)}-(a_{20})^{(3)}=0$ and $(b_{21})^{(3)}(u^{(3)})^2+(\tau_2)^{(3)}u^{(3)}-(b_{20})^{(3)}=0$ | 320 |
| **Definition of** $(m_1)^{(3)},(m_2)^{(3)},(\mu_1)^{(3)},(\mu_2)^{(3)}$:- | 321 |

If we define $(m_1)^{(3)}, (m_2)^{(3)}, (\mu_1)^{(3)}, (\mu_2)^{(3)}$ by

$(m_2)^{(3)} = (\nu_0)^{(3)}, (m_1)^{(3)} = (\nu_1)^{(3)}$, ***if*** $(\nu_0)^{(3)} < (\nu_1)^{(3)}$

$(m_2)^{(3)} = (\nu_1)^{(3)}, (m_1)^{(3)} = (\bar{\nu}_1)^{(3)}$, ***if*** $(\nu_1)^{(3)} < (\nu_0)^{(3)} < (\bar{\nu}_1)^{(3)}$,

and $\boxed{(\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0}}$

$(m_2)^{(3)} = (\nu_1)^{(3)}, (m_1)^{(3)} = (\nu_0)^{(3)}$, ***if*** $(\bar{\nu}_1)^{(3)} < (\nu_0)^{(3)}$

and analogously

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$(\mu_2)^{(3)} = (u_0)^{(3)}, (\mu_1)^{(3)} = (u_1)^{(3)}$, ***if*** $(u_0)^{(3)} < (u_1)^{(3)}$

$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (\bar{u}_1)^{(3)}$, ***if*** $(u_1)^{(3)} < (u_0)^{(3)} < (\bar{u}_1)^{(3)}$, and $\boxed{(u_0)^{(3)} = \frac{T_{20}^0}{T_{21}^0}}$

$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (u_0)^{(3)}$, ***if*** $(\bar{u}_1)^{(3)} < (u_0)^{(3)}$

Then the solution of global equations satisfies the inequalities

$$G_{20}^0 e^{((S_1)^{(3)} - (p_{20})^{(3)})t} \le G_{20}(t) \le G_{20}^0 e^{(S_1)^{(3)} t}$$

$(p_i)^{(3)}$ is defined by equation

$$\frac{1}{(m_1)^{(3)}} G_{20}^0 e^{((S_1)^{(3)} - (p_{20})^{(3)})t} \le G_{21}(t) \le \frac{1}{(m_2)^{(3)}} G_{20}^0 e^{(S_1)^{(3)} t} \quad (323)$$

$$\left(\frac{(a_{22})^{(3)} G_{20}^0}{(m_1)^{(3)}((S_1)^{(3)} - (p_{20})^{(3)} - (S_2)^{(3)})}\left[e^{((S_1)^{(3)} - (p_{20})^{(3)})t} - e^{-(S_2)^{(3)} t}\right] + G_{22}^0 e^{-(S_2)^{(3)} t} \le G_{22}(t) \le \right.$$

$$\left. \frac{(a_{22})^{(3)} G_{20}^0}{(m_2)^{(3)}((S_1)^{(3)} - (a_{22}')^{(3)})}\left[e^{(S_1)^{(3)} t} - e^{-(a_{22}')^{(3)} t}\right] + G_{22}^0 e^{-(a_{22}')^{(3)} t}\right) \quad (324)$$

$$\boxed{T_{20}^0 e^{(R_1)^{(3)} t} \le T_{20}(t) \le T_{20}^0 e^{((R_1)^{(3)} + (r_{20})^{(3)})t}} \quad (325)$$

$$\frac{1}{(\mu_1)^{(3)}} T_{20}^0 e^{(R_1)^{(3)} t} \le T_{20}(t) \le \frac{1}{(\mu_2)^{(3)}} T_{20}^0 e^{((R_1)^{(3)} + (r_{20})^{(3)})t} \quad (326)$$

$$\frac{(b_{22})^{(3)} T_{20}^0}{(\mu_1)^{(3)}((R_1)^{(3)} - (b_{22}')^{(3)})}\left[e^{(R_1)^{(3)} t} - e^{-(b_{22}')^{(3)} t}\right] + T_{22}^0 e^{-(b_{22}')^{(3)} t} \le T_{22}(t) \le$$

$$\frac{(a_{22})^{(3)} T_{20}^0}{(\mu_2)^{(3)}((R_1)^{(3)} + (r_{20})^{(3)} + (R_2)^{(3)})}\left[e^{((R_1)^{(3)} + (r_{20})^{(3)})t} - e^{-(R_2)^{(3)} t}\right] + T_{22}^0 e^{-(R_2)^{(3)} t} \quad (327)$$

**Definition of** $(S_1)^{(3)}, (S_2)^{(3)}, (R_1)^{(3)}, (R_2)^{(3)}$:- (328)

Where $(S_1)^{(3)} = (a_{20})^{(3)}(m_2)^{(3)} - (a_{20}')^{(3)}$

$(S_2)^{(3)} = (a_{22})^{(3)} - (p_{22})^{(3)}$

$(R_1)^{(3)} = (b_{20})^{(3)}(\mu_2)^{(3)} - (b'_{20})^{(3)}$

$(R_2)^{(3)} = (b'_{22})^{(3)} - (r_{22})^{(3)}$

**Behavior of the solutions of equation**

**Theorem:** If we denote and define

**Definition of** $(\sigma_1)^{(4)}, (\sigma_2)^{(4)}, (\tau_1)^{(4)}, (\tau_2)^{(4)}$ :

$(\sigma_1)^{(4)}, (\sigma_2)^{(4)}, (\tau_1)^{(4)}, (\tau_2)^{(4)}$ four constants satisfying

$-(\sigma_2)^{(4)} \le -(a'_{24})^{(4)} + (a'_{25})^{(4)} - (a''_{24})^{(4)}(T_{25}, t) + (a''_{25})^{(4)}(T_{25}, t) \le -(\sigma_1)^{(4)}$

$-(\tau_2)^{(4)} \le -(b'_{24})^{(4)} + (b'_{25})^{(4)} - (b''_{24})^{(4)}((G_{27}), t) - (b''_{25})^{(4)}((G_{27}), t) \le -(\tau_1)^{(4)}$

**Definition of** $(v_1)^{(4)}, (v_2)^{(4)}, (u_1)^{(4)}, (u_2)^{(4)}, v^{(4)}, u^{(4)}$ : 329

By $(v_1)^{(4)} > 0$, $(v_2)^{(4)} < 0$ and respectively $(u_1)^{(4)} > 0$, $(u_2)^{(4)} < 0$ the roots of the equations $(a_{25})^{(4)}\left(v^{(4)}\right)^2 + (\sigma_1)^{(4)}v^{(4)} - (a_{24})^{(4)} = 0$ and $(b_{25})^{(4)}\left(u^{(4)}\right)^2 + (\tau_1)^{(4)}u^{(4)} - (b_{24})^{(4)} = 0$ and

**Definition of** $(\bar{v}_1)^{(4)}, , (\bar{v}_2)^{(4)}, (\bar{u}_1)^{(4)}, (\bar{u}_2)^{(4)}$ : 330

By $(\bar{v}_1)^{(4)} > 0$, $(\bar{v}_2)^{(4)} < 0$ and respectively $(\bar{u}_1)^{(4)} > 0$, $(\bar{u}_2)^{(4)} < 0$ the roots of the equations $(a_{25})^{(4)}\left(v^{(4)}\right)^2 + (\sigma_2)^{(4)}v^{(4)} - (a_{24})^{(4)} = 0$ and $(b_{25})^{(4)}\left(u^{(4)}\right)^2 + (\tau_2)^{(4)}u^{(4)} - (b_{24})^{(4)} = 0$

**Definition of** $(m_1)^{(4)}, (m_2)^{(4)}, (\mu_1)^{(4)}, (\mu_2)^{(4)}, (v_0)^{(4)}$ :-

If we define $(m_1)^{(4)}, (m_2)^{(4)}, (\mu_1)^{(4)}, (\mu_2)^{(4)}$ by

$(m_2)^{(4)} = (v_0)^{(4)}, (m_1)^{(4)} = (v_1)^{(4)}$, ***if*** $(v_0)^{(4)} < (v_1)^{(4)}$

$(m_2)^{(4)} = (v_1)^{(4)}, (m_1)^{(4)} = (\bar{v}_1)^{(4)}$, ***if*** $(v_4)^{(4)} < (v_0)^{(4)} < (\bar{v}_1)^{(4)}$,

and $\boxed{(v_0)^{(4)} = \frac{G^0_{24}}{G^0_{25}}}$

$(m_2)^{(4)} = (v_4)^{(4)}, (m_1)^{(4)} = (v_0)^{(4)}$, ***if*** $(\bar{v}_4)^{(4)} < (v_0)^{(4)}$

and analogously 331

$(\mu_2)^{(4)} = (u_0)^{(4)}, (\mu_1)^{(4)} = (u_1)^{(4)}$, ***if*** $(u_0)^{(4)} < (u_1)^{(4)}$

$(\mu_2)^{(4)} = (u_1)^{(4)}, (\mu_1)^{(4)} = (\bar{u}_1)^{(4)}$, ***if*** $(u_1)^{(4)} < (u_0)^{(4)} < (\bar{u}_1)^{(4)}$,

and $\boxed{(u_0)^{(4)} = \frac{T^0_{24}}{T^0_{25}}}$

$(\mu_2)^{(4)} = (u_1)^{(4)}, (\mu_1)^{(4)} = (u_0)^{(4)}$, ***if*** $(\bar{u}_1)^{(4)} < (u_0)^{(4)}$ where $(u_1)^{(4)}, (\bar{u}_1)^{(4)}$

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| Then the solution of global equations satisfies the inequalities<br>$G_{24}^0 e^{((S_1)^{(4)}-(p_{24})^{(4)})t} \le G_{24}(t) \le G_{24}^0 e^{(S_1)^{(4)}t}$<br>where $(p_i)^{(4)}$ is defined by equation | 332 |
| $\frac{1}{(m_1)^{(4)}} G_{24}^0 e^{((S_1)^{(4)}-(p_{24})^{(4)})t} \le G_{25}(t) \le \frac{1}{(m_2)^{(4)}} G_{24}^0 e^{(S_1)^{(4)}t}$ | 333 |
| $\left(\frac{(a_{26})^{(4)} G_{24}^0}{(m_1)^{(4)}((S_1)^{(4)}-(p_{24})^{(4)}-(S_2)^{(4)})}\left[e^{((S_1)^{(4)}-(p_{24})^{(4)})t} - e^{-(S_2)^{(4)}t}\right] + G_{26}^0 e^{-(S_2)^{(4)}t} \le G_{26}(t) \le$<br>*(a26)4G240(m2)4(S1)4−(a26′)4e(S1)4t−e−(a26′)4t+ G260e−(a26′)4t* | 334 |
| $T_{24}^0 e^{(R_1)^{(4)}t} \le T_{24}(t) \le T_{24}^0 e^{((R_1)^{(4)}+(r_{24})^{(4)})t}$ | |
| $\frac{1}{(\mu_1)^{(4)}} T_{24}^0 e^{(R_1)^{(4)}t} \le T_{24}(t) \le \frac{1}{(\mu_2)^{(4)}} T_{24}^0 e^{((R_1)^{(4)}+(r_{24})^{(4)})t}$ | 335 |
| $\frac{(b_{26})^{(4)} T_{24}^0}{(\mu_1)^{(4)}((R_1)^{(4)}-(b_{26}')^{(4)})}\left[e^{(R_1)^{(4)}t} - e^{-(b_{26}')^{(4)}t}\right] + T_{26}^0 e^{-(b_{26}')^{(4)}t} \le T_{26}(t) \le$<br>$\frac{(a_{26})^{(4)} T_{24}^0}{(\mu_2)^{(4)}((R_1)^{(4)}+(r_{24})^{(4)}+(R_2)^{(4)})}\left[e^{((R_1)^{(4)}+(r_{24})^{(4)})t} - e^{-(R_2)^{(4)}t}\right] + T_{26}^0 e^{-(R_2)^{(4)}t}$ | 336 |
| **Definition of** $(S_1)^{(4)}, (S_2)^{(4)}, (R_1)^{(4)}, (R_2)^{(4)}$:-<br>Where $(S_1)^{(4)} = (a_{24})^{(4)}(m_2)^{(4)} - (a_{24}')^{(4)}$<br>$(S_2)^{(4)} = (a_{26})^{(4)} - (p_{26})^{(4)}$<br>$(R_1)^{(4)} = (b_{24})^{(4)}(\mu_2)^{(4)} - (b_{24}')^{(4)}$<br>$(R_2)^{(4)} = (b_{26}')^{(4)} - (r_{26})^{(4)}$ | 337 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ :<br>$(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ four constants satisfying<br>$-(\sigma_2)^{(5)} \le -(a_{28}')^{(5)} + (a_{29}')^{(5)} - (a_{28}'')^{(5)}(T_{29}, t) + (a_{29}'')^{(5)}(T_{29}, t) \le -(\sigma_1)^{(5)}$<br>$-(\tau_2)^{(5)} \le -(b_{28}')^{(5)} + (b_{29}')^{(5)} - (b_{28}'')^{(5)}((G_{31}), t) - (b_{29}'')^{(5)}((G_{31}), t) \le -(\tau_1)^{(5)}$ | 338 |
| **Definition of** $(\nu_1)^{(5)}, (\nu_2)^{(5)}, (u_1)^{(5)}, (u_2)^{(5)}, \nu^{(5)}, u^{(5)}$ :<br>By $(\nu_1)^{(5)} > 0$ , $(\nu_2)^{(5)} < 0$ and respectively $(u_1)^{(5)} > 0$ , $(u_2)^{(5)} < 0$ the roots of the equations $(a_{29})^{(5)}(\nu^{(5)})^2 + (\sigma_1)^{(5)}\nu^{(5)} - (a_{28})^{(5)} = 0$ | 339 |

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| and $(b_{29})^{(5)}\left(u^{(5)}\right)^2+(\tau_1)^{(5)}u^{(5)}-(b_{28})^{(5)}=0$ and | |
| **Definition of** $(\bar{v}_1)^{(5)},,(\bar{v}_2)^{(5)},(\bar{u}_1)^{(5)},(\bar{u}_2)^{(5)}$ :<br><br>By $(\bar{v}_1)^{(5)}>0$ , $(\bar{v}_2)^{(5)}<0$ and respectively $(\bar{u}_1)^{(5)}>0$ , $(\bar{u}_2)^{(5)}<0$ the roots of the equations $(a_{29})^{(5)}\left(v^{(5)}\right)^2+(\sigma_2)^{(5)}v^{(5)}-(a_{28})^{(5)}=0$ and $(b_{29})^{(5)}\left(u^{(5)}\right)^2+(\tau_2)^{(5)}u^{(5)}-(b_{28})^{(5)}=0$<br>**Definition of** $(m_1)^{(5)}$ , $(m_2)^{(5)}$ , $(\mu_1)^{(5)},(\mu_2)^{(5)},(v_0)^{(5)}$ :-<br><br>If we define $(m_1)^{(5)}$ , $(m_2)^{(5)}$ , $(\mu_1)^{(5)},(\mu_2)^{(5)}$ by<br><br>$(m_2)^{(5)}=(v_0)^{(5)},(m_1)^{(5)}=(v_1)^{(5)}$, ***if*** $(v_0)^{(5)}<(v_1)^{(5)}$<br><br>$(m_2)^{(5)}=(v_1)^{(5)},(m_1)^{(5)}=(\bar{v}_1)^{(5)}$ , ***if*** $(v_1)^{(5)}<(v_0)^{(5)}<(\bar{v}_1)^{(5)}$,<br>and $\boxed{(v_0)^{(5)}=\frac{G_{28}^0}{G_{29}^0}}$<br><br>$(m_2)^{(5)}=(v_1)^{(5)},(m_1)^{(5)}=(v_0)^{(5)}$, ***if*** $(\bar{v}_1)^{(5)}<(v_0)^{(5)}$ | 340 |
| and analogously<br><br>$(\mu_2)^{(5)}=(u_0)^{(5)},(\mu_1)^{(5)}=(u_1)^{(5)}$, ***if*** $(u_0)^{(5)}<(u_1)^{(5)}$<br><br>$(\mu_2)^{(5)}=(u_1)^{(5)},(\mu_1)^{(5)}=(\bar{u}_1)^{(5)}$ , ***if*** $(u_1)^{(5)}<(u_0)^{(5)}<(\bar{u}_1)^{(5)}$,<br>and $\boxed{(u_0)^{(5)}=\frac{T_{28}^0}{T_{29}^0}}$<br><br>$(\mu_2)^{(5)}=(u_1)^{(5)},(\mu_1)^{(5)}=(u_0)^{(5)}$, ***if*** $(\bar{u}_1)^{(5)}<(u_0)^{(5)}$ where $(u_1)^{(5)},(\bar{u}_1)^{(5)}$ | 341 |
| Then the solution of global equations satisfies the inequalities<br><br>$G_{28}^0e^{((S_1)^{(5)}-(p_{28})^{(5)})t}\le G_{28}(t)\le G_{28}^0e^{(S_1)^{(5)}t}$<br><br>where $(p_i)^{(5)}$ is defined by equation | 342 |
| $\frac{1}{(m_5)^{(5)}}G_{28}^0e^{((S_1)^{(5)}-(p_{28})^{(5)})t}\le G_{29}(t)\le\frac{1}{(m_2)^{(5)}}G_{28}^0e^{(S_1)^{(5)}t}$ | 343 |
| $\left(\frac{(a_{30})^{(5)}G_{28}^0}{(m_1)^{(5)}((S_1)^{(5)}-(p_{28})^{(5)}-(S_2)^{(5)})}\left[e^{((S_1)^{(5)}-(p_{28})^{(5)})t}-e^{-(S_2)^{(5)}t}\right]+G_{30}^0e^{-(S_2)^{(5)}t}\le G_{30}(t)\le\right.$<br>*(a30)5G280(m2)5(S1)5−(a30′)5e(S1)5t−e−(a30′)5t+ G300e−(a30′)5t* | 344 |
| $\boxed{T_{28}^0e^{(R_1)^{(5)}t}\le T_{28}(t)\le T_{28}^0e^{((R_1)^{(5)}+(r_{28})^{(5)})t}}$ | 345 |
| $\frac{1}{(\mu_1)^{(5)}}T_{28}^0e^{(R_1)^{(5)}t}\le T_{28}(t)\le\frac{1}{(\mu_2)^{(5)}}T_{28}^0e^{((R_1)^{(5)}+(r_{28})^{(5)})t}$ | 346 |

$\frac{(b_{30})^{(5)}T_{28}^{0}}{(\mu_1)^{(5)}((R_1)^{(5)}-(b_{30}^{'})^{(5)})}\left[e^{(R_1)^{(5)}t}-e^{-(b_{30}^{'})^{(5)}t}\right]+T_{30}^{0}e^{-(b_{30}^{'})^{(5)}t}\leq T_{30}(t)\leq$

$\frac{(a_{30})^{(5)}T_{28}^{0}}{(\mu_2)^{(5)}((R_1)^{(5)}+(r_{28})^{(5)}+(R_2)^{(5)})}\left[e^{((R_1)^{(5)}+(r_{28})^{(5)})t}-e^{-(R_2)^{(5)}t}\right]+T_{30}^{0}e^{-(R_2)^{(5)}t}$ 347

**Definition of** $(S_1)^{(5)},(S_2)^{(5)},(R_1)^{(5)},(R_2)^{(5)}$:- 348

Where $(S_1)^{(5)}=(a_{28})^{(5)}(m_2)^{(5)}-(a_{28}^{'})^{(5)}$

$(S_2)^{(5)}=(a_{30})^{(5)}-(p_{30})^{(5)}$

$(R_1)^{(5)}=(b_{28})^{(5)}(\mu_2)^{(5)}-(b_{28}^{'})^{(5)}$

$(R_2)^{(5)}=(b_{30}^{'})^{(5)}-(r_{30})^{(5)}$

**Behavior of the solutions of equation** 349

**Theorem 2:** If we denote and define

**Definition of** $(\sigma_1)^{(6)},(\sigma_2)^{(6)},(\tau_1)^{(6)},(\tau_2)^{(6)}$ :

$(\sigma_1)^{(6)},(\sigma_2)^{(6)},(\tau_1)^{(6)},(\tau_2)^{(6)}$ four constants satisfying

$-(\sigma_2)^{(6)}\leq-(a_{32}^{'})^{(6)}+(a_{33}^{'})^{(6)}-(a_{32}^{''})^{(6)}(T_{33},t)+(a_{33}^{''})^{(6)}(T_{33},t)\leq-(\sigma_1)^{(6)}$

$-(\tau_2)^{(6)}\leq-(b_{32}^{'})^{(6)}+(b_{33}^{'})^{(6)}-(b_{32}^{''})^{(6)}\left((G_{35}),t\right)-(b_{33}^{''})^{(6)}\left((G_{35}),t\right)\leq-(\tau_1)^{(6)}$

**Definition of** $(\nu_1)^{(6)},(\nu_2)^{(6)},(u_1)^{(6)},(u_2)^{(6)},\nu^{(6)},u^{(6)}$ : 350

By $(\nu_1)^{(6)}>0$ , $(\nu_2)^{(6)}<0$ and respectively $(u_1)^{(6)}>0$ , $(u_2)^{(6)}<0$ the roots of the equations $(a_{33})^{(6)}\left(\nu^{(6)}\right)^2+(\sigma_1)^{(6)}\nu^{(6)}-(a_{32})^{(6)}=0$ and $(b_{33})^{(6)}\left(u^{(6)}\right)^2+(\tau_1)^{(6)}u^{(6)}-(b_{32})^{(6)}=0$ and

**Definition of** $(\bar{\nu}_1)^{(6)},,(\bar{\nu}_2)^{(6)},(\bar{u}_1)^{(6)},(\bar{u}_2)^{(6)}$ : 351

By $(\bar{\nu}_1)^{(6)}>0$ , $(\bar{\nu}_2)^{(6)}<0$ and respectively $(\bar{u}_1)^{(6)}>0$ , $(\bar{u}_2)^{(6)}<0$ the roots of the equations $(a_{33})^{(6)}\left(\nu^{(6)}\right)^2+(\sigma_2)^{(6)}\nu^{(6)}-(a_{32})^{(6)}=0$ and $(b_{33})^{(6)}\left(u^{(6)}\right)^2+(\tau_2)^{(6)}u^{(6)}-(b_{32})^{(6)}=0$

**Definition of** $(m_1)^{(6)},(m_2)^{(6)},(\mu_1)^{(6)},(\mu_2)^{(6)},(\nu_0)^{(6)}$ :-

If we define $(m_1)^{(6)},(m_2)^{(6)},(\mu_1)^{(6)},(\mu_2)^{(6)}$ by

$(m_2)^{(6)}=(\nu_0)^{(6)},(m_1)^{(6)}=(\nu_1)^{(6)},$ ***if*** $(\nu_0)^{(6)}<(\nu_1)^{(6)}$

$(m_2)^{(6)}=(\nu_1)^{(6)},(m_1)^{(6)}=(\bar{\nu}_6)^{(6)}$ , ***if*** $(\nu_1)^{(6)}<(\nu_0)^{(6)}<(\bar{\nu}_1)^{(6)}$,

and $\boxed{(\nu_0)^{(6)}=\frac{G_{32}^{0}}{G_{33}^{0}}}$

$(m_2)^{(6)}=(\nu_1)^{(6)},(m_1)^{(6)}=(\nu_0)^{(6)},$ ***if*** $(\bar{\nu}_1)^{(6)}<(\nu_0)^{(6)}$

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| and analogously<br>$(\mu_2)^{(6)} = (u_0)^{(6)}, (\mu_1)^{(6)} = (u_1)^{(6)},\ \boldsymbol{if}\ (u_0)^{(6)} < (u_1)^{(6)}$<br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (\bar{u}_1)^{(6)}, \boldsymbol{if}\ (u_1)^{(6)} < (u_0)^{(6)} < (\bar{u}_1)^{(6)},$<br>and $\boxed{(u_0)^{(6)} = \frac{T_{32}^0}{T_{33}^0}}$<br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (u_0)^{(6)}, \boldsymbol{if}\ (\bar{u}_1)^{(6)} < (u_0)^{(6)}$ where $(u_1)^{(6)}, (\bar{u}_1)^{(6)}$ | 352 |
| Then the solution of global equations satisfies the inequalities<br>$G_{32}^0 e^{((S_1)^{(6)} - (p_{32})^{(6)})t} \le G_{32}(t) \le G_{32}^0 e^{(S_1)^{(6)} t}$<br>where $(p_i)^{(6)}$ is defined by equation | 353 |
| $\frac{1}{(m_1)^{(6)}} G_{32}^0 e^{((S_1)^{(6)} - (p_{32})^{(6)})t} \le G_{33}(t) \le \frac{1}{(m_2)^{(6)}} G_{32}^0 e^{(S_1)^{(6)} t}$ | 354 |
| $\left(\frac{(a_{34})^{(6)} G_{32}^0}{(m_1)^{(6)}((S_1)^{(6)} - (p_{32})^{(6)} - (S_2)^{(6)})}\left[e^{((S_1)^{(6)} - (p_{32})^{(6)})t} - e^{-(S_2)^{(6)} t}\right] + G_{34}^0 e^{-(S_2)^{(6)} t} \le G_{34}(t) \le\right.$<br>*(a34)6G320(m2)6(S1)6−(a34′)6e(S1)6t−e−(a34′)6t+ G340e−(a34′)6t* | 355 |
| $\boxed{T_{32}^0 e^{(R_1)^{(6)} t} \le T_{32}(t) \le T_{32}^0 e^{((R_1)^{(6)} + (r_{32})^{(6)})t}}$ | 356 |
| $\frac{1}{(\mu_1)^{(6)}} T_{32}^0 e^{(R_1)^{(6)} t} \le T_{32}(t) \le \frac{1}{(\mu_2)^{(6)}} T_{32}^0 e^{((R_1)^{(6)} + (r_{32})^{(6)})t}$ | 357 |
| $\frac{(b_{34})^{(6)} T_{32}^0}{(\mu_1)^{(6)}((R_1)^{(6)} - (b_{34}')^{(6)})}\left[e^{(R_1)^{(6)} t} - e^{-(b_{34}')^{(6)} t}\right] + T_{34}^0 e^{-(b_{34}')^{(6)} t} \le T_{34}(t) \le$<br>$\frac{(a_{34})^{(6)} T_{32}^0}{(\mu_2)^{(6)}((R_1)^{(6)} + (r_{32})^{(6)} + (R_2)^{(6)})}\left[e^{((R_1)^{(6)} + (r_{32})^{(6)})t} - e^{-(R_2)^{(6)} t}\right] + T_{34}^0 e^{-(R_2)^{(6)} t}$ | 358 |
| **Definition of** $(S_1)^{(6)}, (S_2)^{(6)}, (R_1)^{(6)}, (R_2)^{(6)}$:-<br>Where $(S_1)^{(6)} = (a_{32})^{(6)}(m_2)^{(6)} - (a_{32}')^{(6)}$<br>$(S_2)^{(6)} = (a_{34})^{(6)} - (p_{34})^{(6)}$<br>$(R_1)^{(6)} = (b_{32})^{(6)}(\mu_2)^{(6)} - (b_{32}')^{(6)}$<br>$(R_2)^{(6)} = (b_{34}')^{(6)} - (r_{34})^{(6)}$ | 359 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(7)}, (\sigma_2)^{(7)}, (\tau_1)^{(7)}, (\tau_2)^{(7)}$ : | |

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| $(\sigma_1)^{(7)}, (\sigma_2)^{(7)}, (\tau_1)^{(7)}, (\tau_2)^{(7)}$ four constants satisfying<br><br>$-(\sigma_2)^{(7)} \le -(a'_{36})^{(7)} + (a'_{37})^{(7)} - (a''_{36})^{(7)}(T_{37}, t) + (a''_{37})^{(7)}(T_{37}, t) \le -(\sigma_1)^{(7)}$<br><br>$-(\tau_2)^{(7)} \le -(b'_{36})^{(7)} + (b'_{37})^{(7)} - (b''_{36})^{(7)}((G_{39}), t) - (b''_{37})^{(7)}((G_{39}), t) \le -(\tau_1)^{(7)}$ | |
| **Definition of** $(\nu_1)^{(7)}, (\nu_2)^{(7)}, (u_1)^{(7)}, (u_2)^{(7)}, \nu^{(7)}, u^{(7)}$ :<br><br>By $(\nu_1)^{(7)} > 0, (\nu_2)^{(7)} < 0$ and respectively $(u_1)^{(7)} > 0, (u_2)^{(7)} < 0$ the roots of the equations $(a_{37})^{(7)}(\nu^{(7)})^2 + (\sigma_1)^{(7)}\nu^{(7)} - (a_{36})^{(7)} = 0$<br>and $(b_{37})^{(7)}(u^{(7)})^2 + (\tau_1)^{(7)}u^{(7)} - (b_{36})^{(7)} = 0$ and | 361 |
| **Definition of** $(\bar{\nu}_1)^{(7)},, (\bar{\nu}_2)^{(7)}, (\bar{u}_1)^{(7)}, (\bar{u}_2)^{(7)}$ :<br><br>By $(\bar{\nu}_1)^{(7)} > 0, (\bar{\nu}_2)^{(7)} < 0$ and respectively $(\bar{u}_1)^{(7)} > 0, (\bar{u}_2)^{(7)} < 0$ the roots of the equations $(a_{37})^{(7)}(\nu^{(7)})^2 + (\sigma_2)^{(7)}\nu^{(7)} - (a_{36})^{(7)} = 0$<br>and $(b_{37})^{(7)}(u^{(7)})^2 + (\tau_2)^{(7)}u^{(7)} - (b_{36})^{(7)} = 0$<br>**Definition of** $(m_1)^{(7)}, (m_2)^{(7)}, (\mu_1)^{(7)}, (\mu_2)^{(7)}, (\nu_0)^{(7)}$ :-<br><br>If we define $(m_1)^{(7)}, (m_2)^{(7)}, (\mu_1)^{(7)}, (\mu_2)^{(7)}$ by<br><br>$(m_2)^{(7)} = (\nu_0)^{(7)}, (m_1)^{(7)} = (\nu_1)^{(7)}$, ***if*** $(\nu_0)^{(7)} < (\nu_1)^{(7)}$<br><br>$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\bar{\nu}_1)^{(7)}$, ***if*** $(\nu_1)^{(7)} < (\nu_0)^{(7)} < (\bar{\nu}_1)^{(7)}$,<br>and $\boxed{(\nu_0)^{(7)} = \frac{G^0_{36}}{G^0_{37}}}$<br><br>$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\nu_0)^{(7)}$, ***if*** $(\bar{\nu}_1)^{(7)} < (\nu_0)^{(7)}$ | 362 |
| and analogously<br><br>$(\mu_2)^{(7)} = (u_0)^{(7)}, (\mu_1)^{(7)} = (u_1)^{(7)}$, ***if*** $(u_0)^{(7)} < (u_1)^{(7)}$<br><br>$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (\bar{u}_1)^{(7)}$, ***if*** $(u_1)^{(7)} < (u_0)^{(7)} < (\bar{u}_1)^{(7)}$,<br>and $\boxed{(u_0)^{(7)} = \frac{T^0_{36}}{T^0_{37}}}$<br><br>$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (u_0)^{(7)}$, ***if*** $(\bar{u}_1)^{(7)} < (u_0)^{(7)}$ where $(u_1)^{(7)}, (\bar{u}_1)^{(7)}$ | 363 |
| Then the solution of global equations satisfies the inequalities | 364 |

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| $G_{36}^{0}e^{((S_1)^{(7)}-(p_{36})^{(7)})t} \le G_{36}(t) \le G_{36}^{0}e^{(S_1)^{(7)}t}$<br><br>where $(p_i)^{(7)}$ is defined by equation | |
| $\frac{1}{(m_7)^{(7)}} G_{36}^{0}e^{((S_1)^{(7)}-(p_{36})^{(7)})t} \le G_{37}(t) \le \frac{1}{(m_2)^{(7)}} G_{36}^{0}e^{(S_1)^{(7)}t}$ | 365 |
| $(\frac{(a_{38})^{(7)}G_{36}^{0}}{(m_1)^{(7)}((S_1)^{(7)}-(p_{36})^{(7)}-(S_2)^{(7)})}\left[e^{((S_1)^{(7)}-(p_{36})^{(7)})t} - e^{-(S_2)^{(7)}t}\right] + G_{38}^{0}e^{-(S_2)^{(7)}t} \le G_{38}(t) \le$<br>$\frac{(a_{38})^{(7)}G_{36}^{0}}{(m_2)^{(7)}((S_1)^{(7)}-(a_{38}^{'})^{(7)})}[e^{(S_1)^{(7)}t} - e^{-(a_{38}^{'})^{(7)}t}] + G_{38}^{0}e^{-(a_{38}^{'})^{(7)}t})$ | 366 |
| $\boxed{T_{36}^{0}e^{(R_1)^{(7)}t} \le T_{36}(t) \le T_{36}^{0}e^{((R_1)^{(7)}+(r_{36})^{(7)})t}}$ | 367 |
| $\frac{1}{(\mu_1)^{(7)}} T_{36}^{0}e^{(R_1)^{(7)}t} \le T_{36}(t) \le \frac{1}{(\mu_2)^{(7)}} T_{36}^{0}e^{((R_1)^{(7)}+(r_{36})^{(7)})t}$ | 368 |
| $\frac{(b_{38})^{(7)}T_{36}^{0}}{(\mu_1)^{(7)}((R_1)^{(7)}-(b_{38}^{'})^{(7)})}\left[e^{(R_1)^{(7)}t} - e^{-(b_{38}^{'})^{(7)}t}\right] + T_{38}^{0}e^{-(b_{38}^{'})^{(7)}t} \le T_{38}(t) \le$<br><br>$\frac{(a_{38})^{(7)}T_{36}^{0}}{(\mu_2)^{(7)}((R_1)^{(7)}+(r_{36})^{(7)}+(R_2)^{(7)})}\left[e^{((R_1)^{(7)}+(r_{36})^{(7)})t} - e^{-(R_2)^{(7)}t}\right] + T_{38}^{0}e^{-(R_2)^{(7)}t}$ | 369 |
| **Definition of** $(S_1)^{(7)}, (S_2)^{(7)}, (R_1)^{(7)}, (R_2)^{(7)}$:-<br><br>Where $(S_1)^{(7)} = (a_{36})^{(7)}(m_2)^{(7)} - (a_{36}^{'})^{(7)}$<br><br>$(S_2)^{(7)} = (a_{38})^{(7)} - (p_{38})^{(7)}$<br><br>$(R_1)^{(7)} = (b_{36})^{(7)}(\mu_2)^{(7)} - (b_{36}^{'})^{(7)}$<br><br>$(R_2)^{(7)} = (b_{38}^{'})^{(7)} - (r_{38})^{(7)}$ | 370 |
| **Behavior of the solutions of equation**<br><br>**Theorem 2:** If we denote and define<br><br>**Definition of** $(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ :<br><br>$(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ four constants satisfying<br><br>$-(\sigma_2)^{(8)} \le -(a_{40}^{'})^{(8)} + (a_{41}^{'})^{(8)} - (a_{40}^{''})^{(8)}(T_{41}, t) + (a_{41}^{''})^{(8)}(T_{41}, t) \le -(\sigma_1)^{(8)}$<br><br>$-(\tau_2)^{(8)} \le -(b_{40}^{'})^{(8)} + (b_{41}^{'})^{(8)} - (b_{40}^{''})^{(8)}((G_{43}), t) - (b_{41}^{''})^{(8)}((G_{43}), t) \le -(\tau_1)^{(8)}$ | 371 |
| **Definition of** $(v_1)^{(8)}, (v_2)^{(8)}, (u_1)^{(8)}, (u_2)^{(8)}, v^{(8)}, u^{(8)}$ : | 372 |

| | |
|---|---|
| By $(v_1)^{(8)} > 0\,, (v_2)^{(8)} < 0$ and respectively $(u_1)^{(8)} > 0\,, (u_2)^{(8)} < 0$ the roots of the equations $(a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_1)^{(8)}v^{(8)} - (a_{40})^{(8)} = 0$ and $(b_{41})^{(8)}\left(u^{(8)}\right)^2 + (\tau_1)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$ and | |
| **Definition of** $(\bar{v}_1)^{(8)},, (\bar{v}_2)^{(8)}, (\bar{u}_1)^{(8)}, (\bar{u}_2)^{(8)}$ :<br>By $(\bar{v}_1)^{(8)} > 0\,, (\bar{v}_2)^{(8)} < 0$ and respectively $(\bar{u}_1)^{(8)} > 0\,, (\bar{u}_2)^{(8)} < 0$ the<br>roots of the equations $(a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_2)^{(8)}v^{(8)} - (a_{40})^{(8)} = 0$<br>and $(b_{41})^{(8)}\left(u^{(8)}\right)^2 + (\tau_2)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$<br>**Definition of** $(m_1)^{(8)}\,, (m_2)^{(8)}\,, (\mu_1)^{(8)}, (\mu_2)^{(8)}, (v_0)^{(8)}$ :-<br>If we define $(m_1)^{(8)}\,, (m_2)^{(8)}\,, (\mu_1)^{(8)}, (\mu_2)^{(8)}$ by<br>$(m_2)^{(8)} = (v_0)^{(8)}, (m_1)^{(8)} = (v_1)^{(8)},$ ***if*** $(v_0)^{(8)} < (v_1)^{(8)}$<br>$(m_2)^{(8)} = (v_1)^{(8)}, (m_1)^{(8)} = (\bar{v}_1)^{(8)}\,,$ ***if*** $(v_1)^{(8)} < (v_0)^{(8)} < (\bar{v}_1)^{(8)},$<br>and $\boxed{(v_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}}$<br>$(m_2)^{(8)} = (v_1)^{(8)}, (m_1)^{(8)} = (v_0)^{(8)},$ ***if*** $(\bar{v}_1)^{(8)} < (v_0)^{(8)}$ | |
| and analogously<br>$(\mu_2)^{(8)} = (u_0)^{(8)}, (\mu_1)^{(8)} = (u_1)^{(8)},$ ***if*** $(u_0)^{(8)} < (u_1)^{(8)}$<br>$(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (\bar{u}_1)^{(8)}\,,$ ***if*** $(u_1)^{(8)} < (u_0)^{(8)} < (\bar{u}_1)^{(8)},$<br>and $\boxed{(u_0)^{(8)} = \frac{T_{40}^0}{T_{41}^0}}$<br>$(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (u_0)^{(8)},$ ***if*** $(\bar{u}_1)^{(8)} < (u_0)^{(8)}$ where $(u_1)^{(8)}, (\bar{u}_1)^{(8)}$ | 374 |
| Then the solution of global equations satisfies the inequalities<br>$G_{40}^0 e^{((S_1)^{(8)} - (p_{40})^{(8)})t} \le G_{40}(t) \le G_{40}^0 e^{(S_1)^{(8)}t}$<br>where $(p_i)^{(8)}$ is defined by equation | 375 |
| $\frac{1}{(m_1)^{(8)}} G_{40}^0 e^{((S_1)^{(8)} - (p_{40})^{(8)})t} \le G_{41}(t) \le \frac{1}{(m_2)^{(8)}} G_{40}^0 e^{(S_1)^{(8)}t}$ | 376 |

| | |
|---|---|
| $\left(\frac{(a_{42})^{(8)}G_{40}^{0}}{(m_1)^{(8)}((S_1)^{(8)}-(p_{40})^{(8)}-(S_2)^{(8)})}\left[e^{((S_1)^{(8)}-(p_{40})^{(8)})t}-e^{-(S_2)^{(8)}t}\right]+G_{42}^{0}e^{-(S_2)^{(8)}t}\leq G_{42}(t)\leq \frac{(a_{42})^{(8)}G_{40}^{0}}{(m_2)^{(8)}((S_1)^{(8)}-(a_{42}^{'})^{(8)})}[e^{(S_1)^{(8)}t}-e^{-(a_{42}^{'})^{(8)}t}]+G_{42}^{0}e^{-(a_{42}^{'})^{(8)}t}\right)$ | 377 |
| $\boxed{T_{40}^{0}e^{(R_1)^{(8)}t}\leq T_{40}(t)\leq T_{40}^{0}e^{((R_1)^{(8)}+(r_{40})^{(8)})t}}$ | 378 |
| $\frac{1}{(\mu_1)^{(8)}}T_{40}^{0}e^{(R_1)^{(8)}t}\leq T_{40}(t)\leq \frac{1}{(\mu_2)^{(8)}}T_{40}^{0}e^{((R_1)^{(8)}+(r_{40})^{(8)})t}$ | 379 |
| $\frac{(b_{42})^{(8)}T_{40}^{0}}{(\mu_1)^{(8)}((R_1)^{(8)}-(b_{42}^{'})^{(8)})}\left[e^{(R_1)^{(8)}t}-e^{-(b_{42}^{'})^{(8)}t}\right]+T_{42}^{0}e^{-(b_{42}^{'})^{(8)}t}\leq T_{42}(t)\leq \frac{(a_{42})^{(8)}T_{40}^{0}}{(\mu_2)^{(8)}((R_1)^{(8)}+(r_{40})^{(8)}+(R_2)^{(8)})}\left[e^{((R_1)^{(8)}+(r_{40})^{(8)})t}-e^{-(R_2)^{(8)}t}\right]+T_{42}^{0}e^{-(R_2)^{(8)}t}$ | 380 |
| **Definition of** $(S_1)^{(8)},(S_2)^{(8)},(R_1)^{(8)},(R_2)^{(8)}$:-<br>Where $(S_1)^{(8)}=(a_{40})^{(8)}(m_2)^{(8)}-(a_{40}^{'})^{(8)}$<br>$(S_2)^{(8)}=(a_{42})^{(8)}-(p_{42})^{(8)}$<br>$(R_1)^{(8)}=(b_{40})^{(8)}(\mu_2)^{(8)}-(b_{40}^{'})^{(8)}$<br>$(R_2)^{(8)}=(b_{42}^{'})^{(8)}-(r_{42})^{(8)}$ | 381 |
| **Behavior of the solutions of equation 37 to 92**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(9)},(\sigma_2)^{(9)},(\tau_1)^{(9)},(\tau_2)^{(9)}$ :<br>$(\sigma_1)^{(9)},(\sigma_2)^{(9)},(\tau_1)^{(9)},(\tau_2)^{(9)}$ four constants satisfying<br>$-(\sigma_2)^{(9)}\leq -(a_{44}^{'})^{(9)}+(a_{45}^{'})^{(9)}-(a_{44}^{''})^{(9)}(T_{45},t)+(a_{45}^{''})^{(9)}(T_{45},t)\leq -(\sigma_1)^{(9)}$<br>$-(\tau_2)^{(9)}\leq -(b_{44}^{'})^{(9)}+(b_{45}^{'})^{(9)}-(b_{44}^{''})^{(9)}((G_{47}),t)-(b_{45}^{''})^{(9)}((G_{47}),t)\leq -(\tau_1)^{(9)}$ | 382 |
| **Definition of** $(\nu_1)^{(9)},(\nu_2)^{(9)},(u_1)^{(9)},(u_2)^{(9)},\nu^{(9)},u^{(9)}$ :<br>By $(\nu_1)^{(9)}>0$ , $(\nu_2)^{(9)}<0$ and respectively $(u_1)^{(9)}>0$ , $(u_2)^{(9)}<0$ the roots of the equations $(a_{45})^{(9)}(\nu^{(9)})^2+(\sigma_1)^{(9)}\nu^{(9)}-(a_{44})^{(9)}=0$<br>and $(b_{45})^{(9)}(u^{(9)})^2+(\tau_1)^{(9)}u^{(9)}-(b_{44})^{(9)}=0$ and | |
| **Definition of** $(\bar{\nu}_1)^{(9)},,(\bar{\nu}_2)^{(9)},(\bar{u}_1)^{(9)},(\bar{u}_2)^{(9)}$ :<br>By $(\bar{\nu}_1)^{(9)}>0$ , $(\bar{\nu}_2)^{(9)}<0$ and respectively $(\bar{u}_1)^{(9)}>0$ , $(\bar{u}_2)^{(9)}<0$ the roots of the equations $(a_{45})^{(9)}(\nu^{(9)})^2+(\sigma_2)^{(9)}\nu^{(9)}-(a_{44})^{(9)}=0$ | |

and $(b_{45})^{(9)}\left(u^{(9)}\right)^2+(\tau_2)^{(9)}u^{(9)}-(b_{44})^{(9)}=0$

**Definition of** $(m_1)^{(9)},(m_2)^{(9)},(\mu_1)^{(9)},(\mu_2)^{(9)},(\nu_0)^{(9)}$ :-

If we define $(m_1)^{(9)},(m_2)^{(9)},(\mu_1)^{(9)},(\mu_2)^{(9)}$ by

$(m_2)^{(9)}=(\nu_0)^{(9)},(m_1)^{(9)}=(\nu_1)^{(9)},$ ***if*** $(\nu_0)^{(9)}<(\nu_1)^{(9)}$

$(m_2)^{(9)}=(\nu_1)^{(9)},(m_1)^{(9)}=(\bar{\nu}_1)^{(9)}$ , ***if*** $(\nu_1)^{(9)}<(\nu_0)^{(9)}<(\bar{\nu}_1)^{(9)}$,

and $\boxed{(\nu_0)^{(9)}=\frac{G_{44}^0}{G_{45}^0}}$

$(m_2)^{(9)}=(\nu_1)^{(9)},(m_1)^{(9)}=(\nu_0)^{(9)},$ ***if*** $(\bar{\nu}_1)^{(9)}<(\nu_0)^{(9)}$

and analogously

$(\mu_2)^{(9)}=(u_0)^{(9)},(\mu_1)^{(9)}=(u_1)^{(9)},$ ***if*** $(u_0)^{(9)}<(u_1)^{(9)}$

$(\mu_2)^{(9)}=(u_1)^{(9)},(\mu_1)^{(9)}=(\bar{u}_1)^{(9)}$ , ***if*** $(u_1)^{(9)}<(u_0)^{(9)}<(\bar{u}_1)^{(9)}$,

and $\boxed{(u_0)^{(9)}=\frac{T_{44}^0}{T_{45}^0}}$

$(\mu_2)^{(9)}=(u_1)^{(9)},(\mu_1)^{(9)}=(u_0)^{(9)},$ ***if*** $(\bar{u}_1)^{(9)}<(u_0)^{(9)}$ where $(u_1)^{(9)},(\bar{u}_1)^{(9)}$ are defined by 59 and 69 respectively

Then the solution of 19,20,21,22,23 and 24 satisfies the inequalities

$$G_{44}^0e^{((S_1)^{(9)}-(p_{44})^{(9)})t}\le G_{44}(t)\le G_{44}^0e^{(S_1)^{(9)}t}$$

where $(p_i)^{(9)}$ is defined by equation 45

$$\frac{1}{(m_9)^{(9)}}G_{44}^0e^{((S_1)^{(9)}-(p_{44})^{(9)})t}\le G_{45}(t)\le\frac{1}{(m_2)^{(9)}}G_{44}^0e^{(S_1)^{(9)}t}$$

$$\left(\frac{(a_{46})^{(9)}G_{44}^0}{(m_1)^{(9)}((S_1)^{(9)}-(p_{44})^{(9)}-(S_2)^{(9)})}\left[e^{((S_1)^{(9)}-(p_{44})^{(9)})t}-e^{-(S_2)^{(9)}t}\right]+G_{46}^0e^{-(S_2)^{(9)}t}\le G_{46}(t)\le\right.$$

$$\left.\frac{(a_{46})^{(9)}G_{44}^0}{(m_2)^{(9)}((S_1)^{(9)}-(a_{46}')^{(9)})}[e^{(S_1)^{(9)}t}-e^{-(a_{46}')^{(9)}t}]+\ G_{46}^0e^{-(a_{46}')^{(9)}t}\right)$$

$$\boxed{T_{44}^0e^{(R_1)^{(9)}t}\le T_{44}(t)\le T_{44}^0e^{((R_1)^{(9)}+(r_{44})^{(9)})t}}$$

$$\frac{1}{(\mu_1)^{(9)}}T_{44}^0e^{(R_1)^{(9)}t}\le T_{44}(t)\le\frac{1}{(\mu_2)^{(9)}}T_{44}^0e^{((R_1)^{(9)}+(r_{44})^{(9)})t}$$

$$\frac{(b_{46})^{(9)}T_{44}^0}{(\mu_1)^{(9)}((R_1)^{(9)}-(b_{46}')^{(9)})}\left[e^{(R_1)^{(9)}t}-e^{-(b_{46}')^{(9)}t}\right]+T_{46}^0e^{-(b_{46}')^{(9)}t}\le T_{46}(t)\le$$

$$\frac{(a_{46})^{(9)}T_{44}^0}{(\mu_2)^{(9)}((R_1)^{(9)}+(r_{44})^{(9)}+(R_2)^{(9)})}\left[e^{((R_1)^{(9)}+(r_{44})^{(9)})t}-e^{-(R_2)^{(9)}t}\right]+T_{46}^0e^{-(R_2)^{(9)}t}$$

**Definition of** $(S_1)^{(9)}, (S_2)^{(9)}, (R_1)^{(9)}, (R_2)^{(9)}$:-

Where $(S_1)^{(9)} = (a_{44})^{(9)}(m_2)^{(9)} - (a'_{44})^{(9)}$

$(S_2)^{(9)} = (a_{46})^{(9)} - (p_{46})^{(9)}$

$(R_1)^{(9)} = (b_{44})^{(9)}(\mu_2)^{(9)} - (b'_{44})^{(9)}$

$(R_2)^{(9)} = (b'_{46})^{(9)} - (r_{46})^{(9)}$

**Proof :** From global equations we obtain (383)

$$\frac{dv^{(1)}}{dt} = (a_{13})^{(1)} - \left((a'_{13})^{(1)} - (a'_{14})^{(1)} + (a''_{13})^{(1)}(T_{14}, t)\right) - (a''_{14})^{(1)}(T_{14}, t)v^{(1)} - (a_{14})^{(1)}v^{(1)}$$

**Definition of** $v^{(1)}$ :- $\boxed{v^{(1)} = \frac{G_{13}}{G_{14}}}$

It follows

$$-\left((a_{14})^{(1)}\left(v^{(1)}\right)^2 + (\sigma_2)^{(1)}v^{(1)} - (a_{13})^{(1)}\right) \le \frac{dv^{(1)}}{dt} \le -\left((a_{14})^{(1)}\left(v^{(1)}\right)^2 + (\sigma_1)^{(1)}v^{(1)} - (a_{13})^{(1)}\right)$$

From which one obtains

**Definition of** $(\bar{v}_1)^{(1)}, (v_0)^{(1)}$ :-

For $0 < \boxed{(v_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}} < (v_1)^{(1)} < (\bar{v}_1)^{(1)}$

$$v^{(1)}(t) \ge \frac{(v_1)^{(1)} + (C)^{(1)}(v_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((v_1)^{(1)} - (v_0)^{(1)}\right)t\right]}}{1 + (C)^{(1)}e^{\left[-(a_{14})^{(1)}\left((v_1)^{(1)} - (v_0)^{(1)}\right)t\right]}} \quad , \quad \boxed{(C)^{(1)} = \frac{(v_1)^{(1)} - (v_0)^{(1)}}{(v_0)^{(1)} - (v_2)^{(1)}}}$$

it follows $(v_0)^{(1)} \le v^{(1)}(t) \le (v_1)^{(1)}$

In the same manner , we get (384)

$$v^{(1)}(t) \le \frac{(\bar{v}_1)^{(1)} + (\bar{C})^{(1)}(\bar{v}_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{v}_1)^{(1)} - (\bar{v}_2)^{(1)}\right)t\right]}}{1 + (\bar{C})^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{v}_1)^{(1)} - (\bar{v}_2)^{(1)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(1)} = \frac{(\bar{v}_1)^{(1)} - (v_0)^{(1)}}{(v_0)^{(1)} - (\bar{v}_2)^{(1)}}}$$

From which we deduce $(v_0)^{(1)} \le v^{(1)}(t) \le (\bar{v}_1)^{(1)}$

If $0 < (v_1)^{(1)} < (v_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0} < (\bar{v}_1)^{(1)}$ we find like in the previous case, (385)

$$(v_1)^{(1)} \le \frac{(v_1)^{(1)} + (C)^{(1)}(v_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((v_1)^{(1)} - (v_2)^{(1)}\right)t\right]}}{1 + (C)^{(1)}e^{\left[-(a_{14})^{(1)}\left((v_1)^{(1)} - (v_2)^{(1)}\right)t\right]}} \le v^{(1)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(1)}+(\bar{C})^{(1)}(\bar{\nu}_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}}{1+(\bar{C})^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}} \leq (\bar{\nu}_1)^{(1)}$$

If $0 < (\nu_1)^{(1)} \leq (\bar{\nu}_1)^{(1)} \leq \boxed{(\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}}$ , we obtain 386

$$(\nu_1)^{(1)} \leq \nu^{(1)}(t) \leq \frac{(\bar{\nu}_1)^{(1)}+(\bar{C})^{(1)}(\bar{\nu}_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}}{1+(\bar{C})^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}} \leq (\nu_0)^{(1)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(1)}(t)$ :-

$$(m_2)^{(1)} \leq \nu^{(1)}(t) \leq (m_1)^{(1)}, \quad \boxed{\nu^{(1)}(t) = \frac{G_{13}(t)}{G_{14}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(1)}(t)$ :-

$$(\mu_2)^{(1)} \leq u^{(1)}(t) \leq (\mu_1)^{(1)}, \quad \boxed{u^{(1)}(t) = \frac{T_{13}(t)}{T_{14}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{13}'')^{(1)} = (a_{14}'')^{(1)}, then\ (\sigma_1)^{(1)} = (\sigma_2)^{(1)}$ and in this case $(\nu_1)^{(1)} = (\bar{\nu}_1)^{(1)}$ if in addition $(\nu_0)^{(1)} = (\nu_1)^{(1)}$ then $\nu^{(1)}(t) = (\nu_0)^{(1)}$ and as a consequence $G_{13}(t) = (\nu_0)^{(1)}G_{14}(t)$ this also defines $(\nu_0)^{(1)}$ for the special case

Analogously if $(b_{13}'')^{(1)} = (b_{14}'')^{(1)}, then\ (\tau_1)^{(1)} = (\tau_2)^{(1)}$ and then

$(u_1)^{(1)} = (\bar{u}_1)^{(1)}$if in addition $(u_0)^{(1)} = (u_1)^{(1)}$ then $T_{13}(t) = (u_0)^{(1)}T_{14}(t)$ This is an important consequence of the relation between $(\nu_1)^{(1)}$ and $(\bar{\nu}_1)^{(1)}$, and definition of $(u_0)^{(1)}$.

**Proof** : From global equations we obtain 387

$$\frac{d\nu^{(2)}}{dt} = (a_{16})^{(2)} - \left((a_{16}')^{(2)} - (a_{17}')^{(2)} + (a_{16}'')^{(2)}(T_{17}, t)\right) - (a_{17}'')^{(2)}(T_{17}, t)\nu^{(2)} - (a_{17})^{(2)}\nu^{(2)}$$

**Definition of** $\nu^{(2)}$ :- $\boxed{\nu^{(2)} = \frac{G_{16}}{G_{17}}}$ 388

It follows 389

$$-\left((a_{17})^{(2)}\left(\nu^{(2)}\right)^2 + (\sigma_2)^{(2)}\nu^{(2)} - (a_{16})^{(2)}\right) \leq \frac{d\nu^{(2)}}{dt} \leq -\left((a_{17})^{(2)}\left(\nu^{(2)}\right)^2 + (\sigma_1)^{(2)}\nu^{(2)} - (a_{16})^{(2)}\right)$$

From which one obtains 390

**Definition of** $(\bar{\nu}_1)^{(2)}, (\nu_0)^{(2)}$ :-

For $0 < (\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0} < (\nu_1)^{(2)} < (\bar{\nu}_1)^{(2)}$

$$\nu^{(2)}(t) \geq \frac{(\nu_1)^{(2)} + (C)^{(2)}(\nu_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_0)^{(2)}\right)t\right]}}{1 + (C)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_0)^{(2)}\right)t\right]}} \quad , \quad \boxed{(C)^{(2)} = \frac{(\nu_1)^{(2)} - (\nu_0)^{(2)}}{(\nu_0)^{(2)} - (\nu_2)^{(2)}}}$$

it follows $(\nu_0)^{(2)} \leq \nu^{(2)}(t) \leq (\nu_1)^{(2)}$

In the same manner , we get 391

$$\nu^{(2)}(t) \leq \frac{(\bar{\nu}_1)^{(2)} + (\bar{C})^{(2)}(\bar{\nu}_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}}{1 + (\bar{C})^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(2)} = \frac{(\bar{\nu}_1)^{(2)} - (\nu_0)^{(2)}}{(\nu_0)^{(2)} - (\bar{\nu}_2)^{(2)}}}$$

From which we deduce $(\nu_0)^{(2)} \leq \nu^{(2)}(t) \leq (\bar{\nu}_1)^{(2)}$ 392

If $0 < (\nu_1)^{(2)} < (\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0} < (\bar{\nu}_1)^{(2)}$ we find like in the previous case, 393

$$(\nu_1)^{(2)} \leq \frac{(\nu_1)^{(2)} + (C)^{(2)}(\nu_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_2)^{(2)}\right)t\right]}}{1 + (C)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_2)^{(2)}\right)t\right]}} \leq \nu^{(2)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(2)} + (\bar{C})^{(2)}(\bar{\nu}_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}}{1 + (\bar{C})^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}} \leq (\bar{\nu}_1)^{(2)}$$

If $0 < (\nu_1)^{(2)} \leq (\bar{\nu}_1)^{(2)} \leq (\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0}$ , we obtain 394

$$(\nu_1)^{(2)} \leq \nu^{(2)}(t) \leq \frac{(\bar{\nu}_1)^{(2)} + (\bar{C})^{(2)}(\bar{\nu}_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}}{1 + (\bar{C})^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}} \leq (\nu_0)^{(2)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(2)}(t)$ :- 395

$$(m_2)^{(2)} \leq \nu^{(2)}(t) \leq (m_1)^{(2)}, \quad \boxed{\nu^{(2)}(t) = \frac{G_{16}(t)}{G_{17}(t)}}$$

In a completely analogous way, we obtain 396

**Definition of** $u^{(2)}(t)$ :-

$$(\mu_2)^{(2)} \leq u^{(2)}(t) \leq (\mu_1)^{(2)}, \quad \boxed{u^{(2)}(t) = \frac{T_{16}(t)}{T_{17}(t)}}$$

| | |
|---|---|
| Now, using this result and replacing it in global equations we get easily the result stated in the theorem. | |
| **Particular case :**<br><br>If $(a_{16}^{''})^{(2)} = (a_{17}^{''})^{(2)}, then\ (\sigma_1)^{(2)} = (\sigma_2)^{(2)}$ and in this case $(\nu_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ if in addition $(\nu_0)^{(2)} = (\nu_1)^{(2)}$ then $\nu^{(2)}(t) = (\nu_0)^{(2)}$ and as a consequence $G_{16}(t) = (\nu_0)^{(2)} G_{17}(t)$<br><br>Analogously if $(b_{16}^{''})^{(2)} = (b_{17}^{''})^{(2)}, then\ (\tau_1)^{(2)} = (\tau_2)^{(2)}$ and then<br><br>$(u_1)^{(2)} = (\bar{u}_1)^{(2)}$ if in addition $(u_0)^{(2)} = (u_1)^{(2)}$ then $T_{16}(t) = (u_0)^{(2)} T_{17}(t)$ This is an important consequence of the relation between $(\nu_1)^{(2)}$ and $(\bar{\nu}_1)^{(2)}$ | 397 |
| **Proof** : From global equations we obtain<br><br>$\frac{d\nu^{(3)}}{dt} = (a_{20})^{(3)} - \left((a_{20}')^{(3)} - (a_{21}')^{(3)} + (a_{20}^{''})^{(3)}(T_{21}, t)\right) - (a_{21}^{''})^{(3)}(T_{21}, t)\nu^{(3)} - (a_{21})^{(3)}\nu^{(3)}$ | 398 |
| **Definition of** $\nu^{(3)}$ :- $\boxed{\nu^{(3)} = \frac{G_{20}}{G_{21}}}$<br><br>It follows<br><br>$-\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_2)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right) \leq \frac{d\nu^{(3)}}{dt} \leq -\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_1)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right)$ | 399 |
| From which one obtains<br><br>For $0 < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\nu_1)^{(3)} < (\bar{\nu}_1)^{(3)}$<br><br>$\nu^{(3)}(t) \geq \frac{(\nu_1)^{(3)} + (C)^{(3)}(\nu_2)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)} - (\nu_0)^{(3)}\right)t\right]}}{1 + (C)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)} - (\nu_0)^{(3)}\right)t\right]}}$ , $\boxed{(C)^{(3)} = \frac{(\nu_1)^{(3)} - (\nu_0)^{(3)}}{(\nu_0)^{(3)} - (\nu_2)^{(3)}}}$<br><br>it follows $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\nu_1)^{(3)}$ | 400 |
| In the same manner , we get<br><br>$\nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)} + (\bar{C})^{(3)}(\bar{\nu}_2)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)} - (\bar{\nu}_2)^{(3)}\right)t\right]}}{1 + (C)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)} - (\bar{\nu}_2)^{(3)}\right)t\right]}}$ , $\boxed{(\bar{C})^{(3)} = \frac{(\bar{\nu}_1)^{(3)} - (\nu_0)^{(3)}}{(\nu_0)^{(3)} - (\bar{\nu}_2)^{(3)}}}$<br><br>**Definition of** $(\bar{\nu}_1)^{(3)}$ :-<br><br>From which we deduce $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\bar{\nu}_1)^{(3)}$ | 401 |
| If $0 < (\nu_1)^{(3)} < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\bar{\nu}_1)^{(3)}$ we find like in the previous case, | 402 |

$$(\nu_1)^{(3)} \leq \frac{(\nu_1)^{(3)}+(C)^{(3)}(\nu_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}}{1+(C)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}} \leq \nu^{(3)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\bar{\nu}_1)^{(3)}$$

If $0 < (\nu_1)^{(3)} \leq (\bar{\nu}_1)^{(3)} \leq (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0}$ , we obtain

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$$(\nu_1)^{(3)} \leq \nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\nu_0)^{(3)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(3)}(t)$ :-

$$(m_2)^{(3)} \leq \nu^{(3)}(t) \leq (m_1)^{(3)}, \quad \boxed{\nu^{(3)}(t) = \frac{G_{20}(t)}{G_{21}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(3)}(t)$ :-

$$(\mu_2)^{(3)} \leq u^{(3)}(t) \leq (\mu_1)^{(3)}, \quad \boxed{u^{(3)}(t) = \frac{T_{20}(t)}{T_{21}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{20}^{''})^{(3)} = (a_{21}^{''})^{(3)}, then\ (\sigma_1)^{(3)} = (\sigma_2)^{(3)}$ and in this case $(\nu_1)^{(3)} = (\bar{\nu}_1)^{(3)}$ if in addition $(\nu_0)^{(3)} = (\nu_1)^{(3)}$ then $\nu^{(3)}(t) = (\nu_0)^{(3)}$ and as a consequence $G_{20}(t) = (\nu_0)^{(3)}G_{21}(t)$

Analogously if $(b_{20}^{''})^{(3)} = (b_{21}^{''})^{(3)}, then\ (\tau_1)^{(3)} = (\tau_2)^{(3)}$ and then

$(u_1)^{(3)} = (\bar{u}_1)^{(3)}$if in addition $(u_0)^{(3)} = (u_1)^{(3)}$ then $T_{20}(t) = (u_0)^{(3)}T_{21}(t)$ This is an important consequence of the relation between $(\nu_1)^{(3)}$ and $(\bar{\nu}_1)^{(3)}$

**Proof** : From global equations we obtain

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$$\frac{d\nu^{(4)}}{dt} = (a_{24})^{(4)} - \left((a_{24}^{'})^{(4)} - (a_{25}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25},t)\right) - (a_{25}^{''})^{(4)}(T_{25},t)\nu^{(4)} - (a_{25})^{(4)}\nu^{(4)}$$

**Definition of** $\nu^{(4)}$ :- $\quad \boxed{\nu^{(4)} = \frac{G_{24}}{G_{25}}}$

It follows

$$-\left((a_{25})^{(4)}\left(\nu^{(4)}\right)^2 + (\sigma_2)^{(4)}\nu^{(4)} - (a_{24})^{(4)}\right) \leq \frac{d\nu^{(4)}}{dt} \leq -\left((a_{25})^{(4)}\left(\nu^{(4)}\right)^2 + (\sigma_4)^{(4)}\nu^{(4)} - (a_{24})^{(4)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(4)}, (\nu_0)^{(4)}$ :-

For $0 < \boxed{(\nu_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0}} < (\nu_1)^{(4)} < (\bar{\nu}_1)^{(4)}$

$$\nu^{(4)}(t) \geq \frac{(\nu_1)^{(4)} + (C)^{(4)}(\nu_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_0)^{(4)}\right)t\right]}}{4 + (C)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_0)^{(4)}\right)t\right]}} \quad , \quad \boxed{(C)^{(4)} = \frac{(\nu_1)^{(4)} - (\nu_0)^{(4)}}{(\nu_0)^{(4)} - (\nu_2)^{(4)}}}$$

it follows $(\nu_0)^{(4)} \leq \nu^{(4)}(t) \leq (\nu_1)^{(4)}$

In the same manner , we get

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$$\nu^{(4)}(t) \leq \frac{(\bar{\nu}_1)^{(4)} + (\bar{C})^{(4)}(\bar{\nu}_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}}{4 + (\bar{C})^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(4)} = \frac{(\bar{\nu}_1)^{(4)} - (\nu_0)^{(4)}}{(\nu_0)^{(4)} - (\bar{\nu}_2)^{(4)}}}$$

From which we deduce $(\nu_0)^{(4)} \leq \nu^{(4)}(t) \leq (\bar{\nu}_1)^{(4)}$

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If $0 < (\nu_1)^{(4)} < (\nu_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0} < (\bar{\nu}_1)^{(4)}$ we find like in the previous case,

$$(\nu_1)^{(4)} \leq \frac{(\nu_1)^{(4)} + (C)^{(4)}(\nu_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_2)^{(4)}\right)t\right]}}{1 + (C)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_2)^{(4)}\right)t\right]}} \leq \nu^{(4)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(4)} + (\bar{C})^{(4)}(\bar{\nu}_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}}{1 + (\bar{C})^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}} \leq (\bar{\nu}_1)^{(4)}$$

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If $0 < (\nu_1)^{(4)} \leq (\bar{\nu}_1)^{(4)} \leq \boxed{(\nu_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0}}$ , we obtain

$$(\nu_1)^{(4)} \leq \nu^{(4)}(t) \leq \frac{(\bar{\nu}_1)^{(4)} + (\bar{C})^{(4)}(\bar{\nu}_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}}{1 + (\bar{C})^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}} \leq (\nu_0)^{(4)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(4)}(t)$ :-

$(m_2)^{(4)} \leq \nu^{(4)}(t) \leq (m_1)^{(4)}$, $\boxed{\nu^{(4)}(t) = \frac{G_{24}(t)}{G_{25}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(4)}(t)$ :-

$(\mu_2)^{(4)} \leq u^{(4)}(t) \leq (\mu_1)^{(4)}$, $\boxed{u^{(4)}(t) = \frac{T_{24}(t)}{T_{25}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{24}^{''})^{(4)} = (a_{25}^{''})^{(4)}, then\ (\sigma_1)^{(4)} = (\sigma_2)^{(4)}$ and in this case $(\nu_1)^{(4)} = (\bar{\nu}_1)^{(4)}$ if in addition $(\nu_0)^{(4)} =$

$(\nu_1)^{(4)}$ then $\nu^{(4)}(t) = (\nu_0)^{(4)}$ and as a consequence $G_{24}(t) = (\nu_0)^{(4)} G_{25}(t)$ **this also defines** $(\nu_0)^{(4)}$ **for the special case** .

Analogously if $(b_{24}^{''})^{(4)} = (b_{25}^{''})^{(4)}, then\ (\tau_1)^{(4)} = (\tau_2)^{(4)}$ and then $(u_1)^{(4)} = (\bar{u}_4)^{(4)}$ if in addition $(u_0)^{(4)} = (u_1)^{(4)}$ then $T_{24}(t) = (u_0)^{(4)} T_{25}(t)$ This is an important consequence of the relation between $(\nu_1)^{(4)}$ and $(\bar{\nu}_1)^{(4)}$, **and definition of** $(u_0)^{(4)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(5)}}{dt} = (a_{28})^{(5)} - \left((a_{28}^{'})^{(5)} - (a_{29}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}, t)\right) - (a_{29}^{''})^{(5)}(T_{29}, t)\nu^{(5)} - (a_{29})^{(5)}\nu^{(5)}$$

**Definition of** $\nu^{(5)}$ :- $\boxed{\nu^{(5)} = \frac{G_{28}}{G_{29}}}$

It follows

$$-\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_2)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right) \leq \frac{d\nu^{(5)}}{dt} \leq -\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_1)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(5)}, (\nu_0)^{(5)}$ :-

For $0 < \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}} < (\nu_1)^{(5)} < (\bar{\nu}_1)^{(5)}$

$$\nu^{(5)}(t) \geq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}}{5 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}} \quad , \quad \boxed{(C)^{(5)} = \frac{(\nu_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\nu_2)^{(5)}}}$$

it follows $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\nu_1)^{(5)}$

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In the same manner , we get

$$\nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)} + (\bar{C})^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{5 + (\bar{C})^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(5)} = \frac{(\bar{\nu}_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\bar{\nu}_2)^{(5)}}}$$

From which we deduce $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\bar{\nu}_5)^{(5)}$

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If $0 < (\nu_1)^{(5)} < (\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0} < (\bar{\nu}_1)^{(5)}$ we find like in the previous case,

$$(\nu_1)^{(5)} \leq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}}{1 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}} \leq \nu^{(5)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(5)} + (\bar{C})^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{1 + (\bar{C})^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\bar{\nu}_1)^{(5)}$$

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If $0 < (\nu_1)^{(5)} \leq (\bar{\nu}_1)^{(5)} \leq \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}}$ , we obtain

$$(\nu_1)^{(5)} \leq \nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)}+(\bar{C})^{(5)}(\bar{\nu}_2)^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}}{1+(\bar{C})^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\nu_0)^{(5)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(5)}(t)$ :-

$$(m_2)^{(5)} \leq \nu^{(5)}(t) \leq (m_1)^{(5)}, \quad \boxed{\nu^{(5)}(t) = \frac{G_{28}(t)}{G_{29}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(5)}(t)$ :-

$$(\mu_2)^{(5)} \leq u^{(5)}(t) \leq (\mu_1)^{(5)}, \quad \boxed{u^{(5)}(t) = \frac{T_{28}(t)}{T_{29}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{28}^{''})^{(5)} = (a_{29}^{''})^{(5)}, then\ (\sigma_1)^{(5)} = (\sigma_2)^{(5)}$ and in this case $(\nu_1)^{(5)} = (\bar{\nu}_1)^{(5)}$ if in addition $(\nu_0)^{(5)} = (\nu_5)^{(5)}$ then $\nu^{(5)}(t) = (\nu_0)^{(5)}$ and as a consequence $G_{28}(t) = (\nu_0)^{(5)}G_{29}(t)$ **this also defines** $(\nu_0)^{(5)}$ **for the special case** .

Analogously if $(b_{28}^{''})^{(5)} = (b_{29}^{''})^{(5)}, then\ (\tau_1)^{(5)} = (\tau_2)^{(5)}$ and then $(u_1)^{(5)} = (\bar{u}_1)^{(5)}$if in addition $(u_0)^{(5)} = (u_1)^{(5)}$ then $T_{28}(t) = (u_0)^{(5)}T_{29}(t)$ This is an important consequence of the relation between $(\nu_1)^{(5)}$ and $(\bar{\nu}_1)^{(5)}$, **and definition of** $(u_0)^{(5)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(6)}}{dt} = (a_{32})^{(6)} - \left((a_{32}^{'})^{(6)} - (a_{33}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33}, t)\right) - (a_{33}^{''})^{(6)}(T_{33}, t)\nu^{(6)} - (a_{33})^{(6)}\nu^{(6)}$$

**Definition of** $\nu^{(6)}$ :- $\quad \boxed{\nu^{(6)} = \frac{G_{32}}{G_{33}}}$

It follows

$$-\left((a_{33})^{(6)}\left(\nu^{(6)}\right)^2 + (\sigma_2)^{(6)}\nu^{(6)} - (a_{32})^{(6)}\right) \leq \frac{d\nu^{(6)}}{dt} \leq -\left((a_{33})^{(6)}\left(\nu^{(6)}\right)^2 + (\sigma_1)^{(6)}\nu^{(6)} - (a_{32})^{(6)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(6)}, (\nu_0)^{(6)}$ :-

For $0 < \boxed{(\nu_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0}} < (\nu_1)^{(6)} < (\bar{\nu}_1)^{(6)}$

$$\nu^{(6)}(t) \geq \frac{(\nu_1)^{(6)}+(C)^{(6)}(\nu_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_0)^{(6)}\right)t\right]}}{1+(C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_0)^{(6)}\right)t\right]}} \quad , \quad \boxed{(C)^{(6)} = \frac{(\nu_1)^{(6)}-(\nu_0)^{(6)}}{(\nu_0)^{(6)}-(\nu_2)^{(6)}}}$$

it follows $(\nu_0)^{(6)} \leq \nu^{(6)}(t) \leq (\nu_1)^{(6)}$

In the same manner , we get 413

$$\nu^{(6)}(t) \le \frac{(\bar{\nu}_1)^{(6)}+(\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}}{1+(\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(6)} = \frac{(\bar{\nu}_1)^{(6)}-(\nu_0)^{(6)}}{(\nu_0)^{(6)}-(\bar{\nu}_2)^{(6)}}}$$

From which we deduce $(\nu_0)^{(6)} \le \nu^{(6)}(t) \le (\bar{\nu}_1)^{(6)}$

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If $0 < (\nu_1)^{(6)} < (\nu_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0} < (\bar{\nu}_1)^{(6)}$ we find like in the previous case,

$$(\nu_1)^{(6)} \le \frac{(\nu_1)^{(6)}+(C)^{(6)}(\nu_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_2)^{(6)}\right)t\right]}}{1+(C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_2)^{(6)}\right)t\right]}} \le \nu^{(6)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(6)}+(\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}}{1+(\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}} \le (\bar{\nu}_1)^{(6)}$$

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If $0 < (\nu_1)^{(6)} \le (\bar{\nu}_1)^{(6)} \le \boxed{(\nu_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0}}$ , we obtain

$$(\nu_1)^{(6)} \le \nu^{(6)}(t) \le \frac{(\bar{\nu}_1)^{(6)}+(\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}}{1+(\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}} \le (\nu_0)^{(6)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(6)}(t)$ :-

$$(m_2)^{(6)} \le \nu^{(6)}(t) \le (m_1)^{(6)}, \quad \boxed{\nu^{(6)}(t) = \frac{G_{32}(t)}{G_{33}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(6)}(t)$ :-

$$(\mu_2)^{(6)} \le u^{(6)}(t) \le (\mu_1)^{(6)}, \quad \boxed{u^{(6)}(t) = \frac{T_{32}(t)}{T_{33}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{32}^{''})^{(6)} = (a_{33}^{''})^{(6)}, then\ (\sigma_1)^{(6)} = (\sigma_2)^{(6)}$ and in this case $(\nu_1)^{(6)} = (\bar{\nu}_1)^{(6)}$ if in addition $(\nu_0)^{(6)} = (\nu_1)^{(6)}$ then $\nu^{(6)}(t) = (\nu_0)^{(6)}$ and as a consequence $G_{32}(t) = (\nu_0)^{(6)}G_{33}(t)$ **this also defines** $(\nu_0)^{(6)}$ **for the special case** .

Analogously if $(b_{32}^{''})^{(6)} = (b_{33}^{''})^{(6)}, then\ (\tau_1)^{(6)} = (\tau_2)^{(6)}$ and then $(u_1)^{(6)} = (\bar{u}_1)^{(6)}$ if in addition $(u_0)^{(6)} = (u_1)^{(6)}$ then $T_{32}(t) = (u_0)^{(6)}T_{33}(t)$ This is an important consequence of the relation between $(\nu_1)^{(6)}$ and $(\bar{\nu}_1)^{(6)}$, **and definition of** $(u_0)^{(6)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(7)}}{dt} = (a_{36})^{(7)} - \left((a_{36}^{'})^{(7)} - (a_{37}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37},t)\right) - (a_{37}^{''})^{(7)}(T_{37},t)\nu^{(7)} - (a_{37})^{(7)}\nu^{(7)}$$

**Definition of** $\nu^{(7)}$ :- $\boxed{\nu^{(7)} = \frac{G_{36}}{G_{37}}}$

It follows

$$-\left((a_{37})^{(7)}\left(\nu^{(7)}\right)^2 + (\sigma_2)^{(7)}\nu^{(7)} - (a_{36})^{(7)}\right) \leq \frac{d\nu^{(7)}}{dt} \leq -\left((a_{37})^{(7)}\left(\nu^{(7)}\right)^2 + (\sigma_1)^{(7)}\nu^{(7)} - (a_{36})^{(7)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(7)}, (\nu_0)^{(7)}$ :-

For $0 < \boxed{(\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0}} < (\nu_1)^{(7)} < (\bar{\nu}_1)^{(7)}$

$$\nu^{(7)}(t) \geq \frac{(\nu_1)^{(7)} + (C)^{(7)}(\nu_2)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)} - (\nu_0)^{(7)}\right)t\right]}}{1 + (C)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)} - (\nu_0)^{(7)}\right)t\right]}} \quad , \quad \boxed{(C)^{(7)} = \frac{(\nu_1)^{(7)} - (\nu_0)^{(7)}}{(\nu_0)^{(7)} - (\nu_2)^{(7)}}}$$

it follows $(\nu_0)^{(7)} \leq \nu^{(7)}(t) \leq (\nu_1)^{(7)}$

In the same manner , we get 417

$$\nu^{(7)}(t) \leq \frac{(\bar{\nu}_1)^{(7)} + (\bar{C})^{(7)}(\bar{\nu}_2)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)} - (\bar{\nu}_2)^{(7)}\right)t\right]}}{1 + (\bar{C})^{(7)} e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)} - (\bar{\nu}_2)^{(7)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(7)} = \frac{(\bar{\nu}_1)^{(7)} - (\nu_0)^{(7)}}{(\nu_0)^{(7)} - (\bar{\nu}_2)^{(7)}}}$$

From which we deduce $(\nu_0)^{(7)} \leq \nu^{(7)}(t) \leq (\bar{\nu}_1)^{(7)}$

If $0 < (\nu_1)^{(7)} < (\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0} < (\bar{\nu}_1)^{(7)}$ we find like in the previous case, 418

$$(\nu_1)^{(7)} \leq \frac{(\nu_1)^{(7)} + (C)^{(7)}(\nu_2)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)} - (\nu_2)^{(7)}\right)t\right]}}{1 + (C)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)} - (\nu_2)^{(7)}\right)t\right]}} \leq \nu^{(7)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(7)} + (\bar{C})^{(7)}(\bar{\nu}_2)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)} - (\bar{\nu}_2)^{(7)}\right)t\right]}}{1 + (\bar{C})^{(7)} e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)} - (\bar{\nu}_2)^{(7)}\right)t\right]}} \leq (\bar{\nu}_1)^{(7)}$$

If $0 < (\nu_1)^{(7)} \leq (\bar{\nu}_1)^{(7)} \leq \boxed{(\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0}}$ , we obtain 419

$$(\nu_1)^{(7)} \leq \nu^{(7)}(t) \leq \frac{(\bar{\nu}_1)^{(7)} + (\bar{C})^{(7)}(\bar{\nu}_2)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)} - (\bar{\nu}_2)^{(7)}\right)t\right]}}{1 + (\bar{C})^{(7)} e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)} - (\bar{\nu}_2)^{(7)}\right)t\right]}} \leq (\nu_0)^{(7)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(7)}(t)$ :-

$$(m_2)^{(7)} \leq \nu^{(7)}(t) \leq (m_1)^{(7)}, \quad \boxed{\nu^{(7)}(t) = \frac{G_{36}(t)}{G_{37}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(7)}(t)$ :- 420

$(\mu_2)^{(7)} \leq u^{(7)}(t) \leq (\mu_1)^{(7)}$, $\boxed{u^{(7)}(t) = \frac{T_{36}(t)}{T_{37}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{36}^{''})^{(7)} = (a_{37}^{''})^{(7)}, then\ (\sigma_1)^{(7)} = (\sigma_2)^{(7)}$ and in this case $(\nu_1)^{(7)} = (\bar{\nu}_1)^{(7)}$ if in addition $(\nu_0)^{(7)} = (\nu_1)^{(7)}$ then $\nu^{(7)}(t) = (\nu_0)^{(7)}$ and as a consequence $G_{36}(t) = (\nu_0)^{(7)} G_{37}(t)$ **this also defines** $(\nu_0)^{(7)}$ **for the special case** .

Analogously if $(b_{36}^{''})^{(7)} = (b_{37}^{''})^{(7)}, then\ (\tau_1)^{(7)} = (\tau_2)^{(7)}$ and then $(u_1)^{(7)} = (\bar{u}_1)^{(7)}$if in addition $(u_0)^{(7)} = (u_1)^{(7)}$ then $T_{36}(t) = (u_0)^{(7)} T_{37}(t)$ This is an important consequence of the relation between $(\nu_1)^{(7)}$ and $(\bar{\nu}_1)^{(7)}$, **and definition of** $(u_0)^{(7)}$.

**Proof** : From global equations we obtain 421

$$\frac{d\nu^{(8)}}{dt} = (a_{40})^{(8)} - \left((a_{40}^{'})^{(8)} - (a_{41}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41}, t)\right) - (a_{41}^{''})^{(8)}(T_{41}, t)\nu^{(8)} - (a_{41})^{(8)}\nu^{(8)}$$

**Definition of** $\nu^{(8)}$ :- $\boxed{\nu^{(8)} = \frac{G_{40}}{G_{41}}}$

It follows

$$-\left((a_{41})^{(8)}\left(\nu^{(8)}\right)^2 + (\sigma_2)^{(8)}\nu^{(8)} - (a_{40})^{(8)}\right) \leq \frac{d\nu^{(8)}}{dt} \leq -\left((a_{41})^{(8)}\left(\nu^{(8)}\right)^2 + (\sigma_1)^{(8)}\nu^{(8)} - (a_{40})^{(8)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(8)}, (\nu_0)^{(8)}$ :-

For $0 < \boxed{(\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}} < (\nu_1)^{(8)} < (\bar{\nu}_1)^{(8)}$

$$\nu^{(8)}(t) \geq \frac{(\nu_1)^{(8)} + (C)^{(8)}(\nu_2)^{(8)} e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}}{1 + (C)^{(8)} e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}} \quad , \quad \boxed{(C)^{(8)} = \frac{(\nu_1)^{(8)} - (\nu_0)^{(8)}}{(\nu_0)^{(8)} - (\nu_2)^{(8)}}}$$

it follows $(\nu_0)^{(8)} \leq \nu^{(8)}(t) \leq (\nu_1)^{(8)}$

In the same manner , we get 422

$$\nu^{(8)}(t) \le \frac{(\bar{\nu}_1)^{(8)}+(\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}}{1+(\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(8)} = \frac{(\bar{\nu}_1)^{(8)}-(\nu_0)^{(8)}}{(\nu_0)^{(8)}-(\bar{\nu}_2)^{(8)}}}$$

From which we deduce $(\nu_0)^{(8)} \le \nu^{(8)}(t) \le (\bar{\nu}_8)^{(8)}$

If $0 < (\nu_1)^{(8)} < (\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0} < (\bar{\nu}_1)^{(8)}$ we find like in the previous case, 423

$$(\nu_1)^{(8)} \le \frac{(\nu_1)^{(8)}+(C)^{(8)}(\nu_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)}-(\nu_2)^{(8)}\right)t\right]}}{1+(C)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)}-(\nu_2)^{(8)}\right)t\right]}} \le \nu^{(8)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(8)}+(\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}}{1+(\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}} \le (\bar{\nu}_1)^{(8)}$$

If $0 < (\nu_1)^{(8)} \le (\bar{\nu}_1)^{(8)} \le \boxed{(\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}}$, we obtain 424

$$(\nu_1)^{(8)} \le \nu^{(8)}(t) \le \frac{(\bar{\nu}_1)^{(8)}+(\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}}{1+(\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}} \le (\nu_0)^{(8)}$$

And so with the notation of the first part of condition (c) , we have
**Definition of** $\nu^{(8)}(t)$ :-

$$(m_2)^{(8)} \le \nu^{(8)}(t) \le (m_1)^{(8)}, \quad \boxed{\nu^{(8)}(t) = \frac{G_{40}(t)}{G_{41}(t)}}$$

In a completely analogous way, we obtain
**Definition of** $u^{(8)}(t)$ :-

$$(\mu_2)^{(8)} \le u^{(8)}(t) \le (\mu_1)^{(8)}, \quad \boxed{u^{(8)}(t) = \frac{T_{40}(t)}{T_{41}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{40}^{''})^{(8)} = (a_{41}^{''})^{(8)}, then\ (\sigma_1)^{(8)} = (\sigma_2)^{(8)}$ and in this case $(\nu_1)^{(8)} = (\bar{\nu}_1)^{(8)}$ if in addition $(\nu_0)^{(8)} = (\nu_1)^{(8)}$ then $\nu^{(8)}(t) = (\nu_0)^{(8)}$ and as a consequence $G_{40}(t) = (\nu_0)^{(8)}G_{41}(t)$ **this also defines** $(\nu_0)^{(8)}$ **for the special case** .

Analogously if $(b_{40}^{''})^{(8)} = (b_{41}^{''})^{(8)}, then\ (\tau_1)^{(8)} = (\tau_2)^{(8)}$ and then
$(u_1)^{(8)} = (\bar{u}_1)^{(8)}$ if in addition $(u_0)^{(8)} = (u_1)^{(8)}$ then $T_{40}(t) = (u_0)^{(8)}T_{41}(t)$ This is an important consequence of the relation between $(\nu_1)^{(8)}$ and $(\bar{\nu}_1)^{(8)}$, **and definition of** $(u_0)^{(8)}$. 424

**Proof** : From 99,20,44,22,23,44 we obtain 424 A

$$\frac{dv^{(9)}}{dt} = (a_{44})^{(9)} - \left((a'_{44})^{(9)} - (a'_{45})^{(9)} + (a''_{44})^{(9)}(T_{45},t)\right) - (a''_{45})^{(9)}(T_{45},t)v^{(9)} - (a_{45})^{(9)}v^{(9)}$$

**Definition of** $v^{(9)}$ :- $\boxed{v^{(9)} = \frac{G_{44}}{G_{45}}}$

It follows

$$-\left((a_{45})^{(9)}\left(v^{(9)}\right)^2 + (\sigma_2)^{(9)}v^{(9)} - (a_{44})^{(9)}\right) \le \frac{dv^{(9)}}{dt} \le -\left((a_{45})^{(9)}\left(v^{(9)}\right)^2 + (\sigma_1)^{(9)}v^{(9)} - (a_{44})^{(9)}\right)$$

From which one obtains

**Definition of** $(\bar{v}_1)^{(9)}, (v_0)^{(9)}$ :-

For $0 < \boxed{(v_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}}} < (v_1)^{(9)} < (\bar{v}_1)^{(9)}$

$$v^{(9)}(t) \ge \frac{(v_1)^{(9)} + (C)^{(9)}(v_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((v_1)^{(9)} - (v_0)^{(9)}\right)t\right]}}{1 + (C)^{(9)}e^{\left[-(a_{45})^{(9)}\left((v_1)^{(9)} - (v_0)^{(9)}\right)t\right]}} \quad , \quad \boxed{(C)^{(9)} = \frac{(v_1)^{(9)} - (v_0)^{(9)}}{(v_0)^{(9)} - (v_2)^{(9)}}}$$

it follows $(v_0)^{(9)} \le v^{(9)}(t) \le (v_9)^{(9)}$

In the same manner , we get

$$v^{(9)}(t) \le \frac{(\bar{v}_1)^{(9)} + (\bar{C})^{(9)}(\bar{v}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)} - (\bar{v}_2)^{(9)}\right)t\right]}}{1 + (\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)} - (\bar{v}_2)^{(9)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(9)} = \frac{(\bar{v}_1)^{(9)} - (v_0)^{(9)}}{(v_0)^{(9)} - (\bar{v}_2)^{(9)}}}$$

From which we deduce $(v_0)^{(9)} \le v^{(9)}(t) \le (\bar{v}_1)^{(9)}$

If $0 < (v_1)^{(9)} < (v_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}} < (\bar{v}_1)^{(9)}$ we find like in the previous case,

$$(v_1)^{(9)} \le \frac{(v_1)^{(9)} + (C)^{(9)}(v_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((v_1)^{(9)} - (v_2)^{(9)}\right)t\right]}}{1 + (C)^{(9)}e^{\left[-(a_{45})^{(9)}\left((v_1)^{(9)} - (v_2)^{(9)}\right)t\right]}} \le v^{(9)}(t) \le$$

$$\frac{(\bar{v}_1)^{(9)} + (\bar{C})^{(9)}(\bar{v}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)} - (\bar{v}_2)^{(9)}\right)t\right]}}{1 + (\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)} - (\bar{v}_2)^{(9)}\right)t\right]}} \le (\bar{v}_1)^{(9)}$$

If $0 < (v_1)^{(9)} \le (\bar{v}_1)^{(9)} \le \boxed{(v_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}}}$ , we obtain

$$(v_1)^{(9)} \le v^{(9)}(t) \le \frac{(\bar{v}_1)^{(9)} + (\bar{C})^{(9)}(\bar{v}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)} - (\bar{v}_2)^{(9)}\right)t\right]}}{1 + (\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)} - (\bar{v}_2)^{(9)}\right)t\right]}} \le (v_0)^{(9)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $v^{(9)}(t)$ :-

$$(m_2)^{(9)} \le v^{(9)}(t) \le (m_1)^{(9)}, \quad \boxed{v^{(9)}(t) = \frac{G_{44}(t)}{G_{45}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(9)}(t)$ :-

$(\mu_2)^{(9)} \leq u^{(9)}(t) \leq (\mu_1)^{(9)}$, $\boxed{u^{(9)}(t) = \frac{T_{44}(t)}{T_{45}(t)}}$

Now, using this result and replacing it in 99, 20,44,22,23, and 44 we get easily the result stated in the theorem.

**Particular case :**

If $(a_{44}^{''})^{(9)} = (a_{45}^{''})^{(9)}, then\ (\sigma_1)^{(9)} = (\sigma_2)^{(9)}$ and in this case $(\nu_1)^{(9)} = (\bar{\nu}_1)^{(9)}$ if in addition $(\nu_0)^{(9)} = (\nu_1)^{(9)}$ then $\nu^{(9)}(t) = (\nu_0)^{(9)}$ and as a consequence $G_{44}(t) = (\nu_0)^{(9)} G_{45}(t)$ **this also defines** $(\nu_0)^{(9)}$ **for the special case** .

Analogously if $(b_{44}^{''})^{(9)} = (b_{45}^{''})^{(9)}, then\ (\tau_1)^{(9)} = (\tau_2)^{(9)}$ and then $(u_1)^{(9)} = (\bar{u}_1)^{(9)}$if in addition $(u_0)^{(9)} = (u_1)^{(9)}$ then $T_{44}(t) = (u_0)^{(9)} T_{45}(t)$ This is an important consequence of the relation between $(\nu_1)^{(9)}$ and $(\bar{\nu}_1)^{(9)}$, **and definition of** $(u_0)^{(9)}$.

We can prove the following 425

**Theorem :** If $(a_i^{''})^{(1)} and\ (b_i^{''})^{(1)}$ are independent on $t$ , and the conditions with the notations

$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} < 0$

$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a_{13})^{(1)}(p_{13})^{(1)} + (a_{14}^{'})^{(1)}(p_{14})^{(1)} + (p_{13})^{(1)}(p_{14})^{(1)} > 0$

$(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} > 0$ ,

$(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} - (b_{13}^{'})^{(1)}(r_{14})^{(1)} - (b_{14}^{'})^{(1)}(r_{14})^{(1)} + (r_{13})^{(1)}(r_{14})^{(1)} < 0$

$with\ (p_{13})^{(1)}, (r_{14})^{(1)}$ as defined by equation are satisfied , then the system

**Theorem :** If $(a_i^{''})^{(2)} and\ (b_i^{''})^{(2)}$ are independent on t , and the conditions with the notations 426

$(a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} < 0$ 427

$(a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a_{16})^{(2)}(p_{16})^{(2)} + (a_{17}^{'})^{(2)}(p_{17})^{(2)} + (p_{16})^{(2)}(p_{17})^{(2)} > 0$ 428

$(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} > 0$ , 429

$(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} - (b_{16}^{'})^{(2)}(r_{17})^{(2)} - (b_{17}^{'})^{(2)}(r_{17})^{(2)} + (r_{16})^{(2)}(r_{17})^{(2)} < 0$ 430

$with\ (p_{16})^{(2)}, (r_{17})^{(2)}$ as defined by equation are satisfied , then the system

**Theorem :** If $(a_i^{''})^{(3)} and\ (b_i^{''})^{(3)}$ are independent on $t$ , and the conditions with the notations 431

$(a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} < 0$

$(a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a_{20})^{(3)}(p_{20})^{(3)} + (a_{21}^{'})^{(3)}(p_{21})^{(3)} + (p_{20})^{(3)}(p_{21})^{(3)} > 0$

$(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} > 0$ ,

$(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} - (b_{20}^{'})^{(3)}(r_{21})^{(3)} - (b_{21}^{'})^{(3)}(r_{21})^{(3)} + (r_{20})^{(3)}(r_{21})^{(3)} < 0$

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| $with\ \ (p_{20})^{(3)}, (r_{21})^{(3)}$ as defined by equation are satisfied , then the system | |
| We can prove the following<br><br>**Theorem :** If $(a_i^{''})^{(4)} and\ (b_i^{''})^{(4)}$ are independent on $t$ , and the conditions with the notations<br><br>$(a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} < 0$<br><br>$(a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a_{24})^{(4)}(p_{24})^{(4)} + (a_{25}^{'})^{(4)}(p_{25})^{(4)} + (p_{24})^{(4)}(p_{25})^{(4)} > 0$<br><br>$(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} > 0$ ,<br><br>$(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} - (b_{24}^{'})^{(4)}(r_{25})^{(4)} - (b_{25}^{'})^{(4)}(r_{25})^{(4)} + (r_{24})^{(4)}(r_{25})^{(4)} < 0$<br><br>$with\ \ (p_{24})^{(4)}, (r_{25})^{(4)}$ as defined by equation are satisfied , then the system | 432 |
| **Theorem :** If $(a_i^{''})^{(5)} and\ (b_i^{''})^{(5)}$ are independent on $t$ , and the conditions with the notations<br><br>$(a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} < 0$<br><br>$(a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a_{28})^{(5)}(p_{28})^{(5)} + (a_{29}^{'})^{(5)}(p_{29})^{(5)} + (p_{28})^{(5)}(p_{29})^{(5)} > 0$<br><br>$(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} > 0$ ,<br><br>$(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} - (b_{28}^{'})^{(5)}(r_{29})^{(5)} - (b_{29}^{'})^{(5)}(r_{29})^{(5)} + (r_{28})^{(5)}(r_{29})^{(5)} < 0$<br><br>$with\ \ (p_{28})^{(5)}, (r_{29})^{(5)}$ as defined by equation are satisfied , then the system | 433 |
| **Theorem** If $(a_i^{''})^{(6)} and\ (b_i^{''})^{(6)}$ are independent on $t$ , and the conditions with the notations<br><br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} < 0$<br><br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32})^{(6)}(p_{32})^{(6)} + (a_{33}^{'})^{(6)}(p_{33})^{(6)} + (p_{32})^{(6)}(p_{33})^{(6)} > 0$<br><br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} > 0$ ,<br><br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} - (b_{32}^{'})^{(6)}(r_{33})^{(6)} - (b_{33}^{'})^{(6)}(r_{33})^{(6)} + (r_{32})^{(6)}(r_{33})^{(6)} < 0$<br><br>$with\ \ (p_{32})^{(6)}, (r_{33})^{(6)}$ as defined by equation are satisfied , then the system | 434 |
| **Theorem :** If $(a_i^{''})^{(7)} and\ (b_i^{''})^{(7)}$ are independent on $t$ , and the conditions with the notations<br><br>$(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} < 0$<br><br>$(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36})^{(7)}(p_{36})^{(7)} + (a_{37}^{'})^{(7)}(p_{37})^{(7)} + (p_{36})^{(7)}(p_{37})^{(7)} > 0$<br><br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} > 0$ ,<br><br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} - (b_{36}^{'})^{(7)}(r_{37})^{(7)} - (b_{37}^{'})^{(7)}(r_{37})^{(7)} + (r_{36})^{(7)}(r_{37})^{(7)} < 0$ | 435 |

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| $with\ \ (p_{36})^{(7)}, (r_{37})^{(7)}$ as defined by equation are satisfied , then the system | |
| **Theorem :** If $(a_i^{''})^{(8)} and\ (b_i^{''})^{(8)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} < 0$<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40})^{(8)}(p_{40})^{(8)} + (a_{41}^{'})^{(8)}(p_{41})^{(8)} + (p_{40})^{(8)}(p_{41})^{(8)} > 0$<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} > 0$ ,<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} - (b_{40}^{'})^{(8)}(r_{41})^{(8)} - (b_{41}^{'})^{(8)}(r_{41})^{(8)} + (r_{40})^{(8)}(r_{41})^{(8)} < 0$<br>$with\ \ (p_{40})^{(8)}, (r_{41})^{(8)}$ as defined by equation are satisfied , then the system | 436 |
| **Theorem :** If $(a_i^{''})^{(9)} and\ (b_i^{''})^{(9)}$ are independent on $t$ , and the conditions (with the notations 45,46,27,28)<br>$(a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} < 0$<br>$(a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44})^{(9)}(p_{44})^{(9)} + (a_{45}^{'})^{(9)}(p_{45})^{(9)} + (p_{44})^{(9)}(p_{45})^{(9)} > 0$<br>$(b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} > 0$ ,<br>$(b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} - (b_{44}^{'})^{(9)}(r_{45})^{(9)} - (b_{45}^{'})^{(9)}(r_{45})^{(9)} + (r_{44})^{(9)}(r_{45})^{(9)} < 0$<br>$with\ \ (p_{44})^{(9)}, (r_{45})^{(9)}$ as defined by equation 45 are satisfied , then the system | 436 A |
| $(a_{13})^{(1)}G_{14} - \left[(a_{13}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14})\right]G_{13} = 0$ | 437 |
| $(a_{14})^{(1)}G_{13} - \left[(a_{14}^{'})^{(1)} + (a_{14}^{''})^{(1)}(T_{14})\right]G_{14} = 0$ | 438 |
| $(a_{15})^{(1)}G_{14} - \left[(a_{15}^{'})^{(1)} + (a_{15}^{''})^{(1)}(T_{14})\right]G_{15} = 0$ | 439 |
| $(b_{13})^{(1)}T_{14} - [(b_{13}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G)\ ]T_{13} = 0$ | 440 |
| $(b_{14})^{(1)}T_{13} - [(b_{14}^{'})^{(1)} - (b_{14}^{''})^{(1)}(G)\ ]T_{14} = 0$ | 441 |
| $(b_{15})^{(1)}T_{14} - [(b_{15}^{'})^{(1)} - (b_{15}^{''})^{(1)}(G)\ ]T_{15} = 0$ | 442 |
| has a unique positive solution , which is an equilibrium solution for the system | |
| $(a_{16})^{(2)}G_{17} - \left[(a_{16}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17})\right]G_{16} = 0$ | 443 |
| $(a_{17})^{(2)}G_{16} - \left[(a_{17}^{'})^{(2)} + (a_{17}^{''})^{(2)}(T_{17})\right]G_{17} = 0$ | 444 |
| $(a_{18})^{(2)}G_{17} - \left[(a_{18}^{'})^{(2)} + (a_{18}^{''})^{(2)}(T_{17})\right]G_{18} = 0$ | 445 |

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| $(b_{16})^{(2)}T_{17} - [(b_{16}^{'})^{(2)} - (b_{16}^{''})^{(2)}(G_{19})\,]T_{16} = 0$ | 446 |
| $(b_{17})^{(2)}T_{16} - [(b_{17}^{'})^{(2)} - (b_{17}^{''})^{(2)}(G_{19})\,]T_{17} = 0$ | 447 |
| $(b_{18})^{(2)}T_{17} - [(b_{18}^{'})^{(2)} - (b_{18}^{''})^{(2)}(G_{19})\,]T_{18} = 0$ | 448 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{20})^{(3)}G_{21} - \left[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21})\right]G_{20} = 0$ | 449 |
| $(a_{21})^{(3)}G_{20} - \left[(a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)}(T_{21})\right]G_{21} = 0$ | 450 |
| $(a_{22})^{(3)}G_{21} - \left[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21})\right]G_{22} = 0$ | 451 |
| $(b_{20})^{(3)}T_{21} - [(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23})\,]T_{20} = 0$ | 452 |
| $(b_{21})^{(3)}T_{20} - [(b_{21}^{'})^{(3)} - (b_{21}^{''})^{(3)}(G_{23})\,]T_{21} = 0$ | 453 |
| $(b_{22})^{(3)}T_{21} - [(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23})\,]T_{22} = 0$ | 454 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25})\right]G_{24} = 0$ | 455 |
| $(a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25})\right]G_{25} = 0$ | 456 |
| $(a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25})\right]G_{26} = 0$ | 457 |
| $(b_{24})^{(4)}T_{25} - [(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}\big((G_{27})\big)\,]T_{24} = 0$ | 458 |
| $(b_{25})^{(4)}T_{24} - [(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}\big((G_{27})\big)\,]T_{25} = 0$ | 459 |
| $(b_{26})^{(4)}T_{25} - [(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}\big((G_{27})\big)\,]T_{26} = 0$ | 460 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29})\right]G_{28} = 0$ | 461 |
| $(a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29})\right]G_{29} = 0$ | 462 |
| $(a_{30})^{(5)}G_{29} - \left[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29})\right]G_{30} = 0$ | 463 |
| $(b_{28})^{(5)}T_{29} - [(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}(G_{31})\,]T_{28} = 0$ | 464 |
| $(b_{29})^{(5)}T_{28} - [(b_{29}^{'})^{(5)} - (b_{29}^{''})^{(5)}(G_{31})\,]T_{29} = 0$ | 465 |

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| $(b_{30})^{(5)}T_{29} - [(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}(G_{31})\,]T_{30} = 0$ | 466 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{32})^{(6)}G_{33} - \left[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33})\right]G_{32} = 0$ | 467 |
| $(a_{33})^{(6)}G_{32} - \left[(a_{33}^{'})^{(6)} + (a_{33}^{''})^{(6)}(T_{33})\right]G_{33} = 0$ | 468 |
| $(a_{34})^{(6)}G_{33} - \left[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33})\right]G_{34} = 0$ | 469 |
| $(b_{32})^{(6)}T_{33} - [(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}(G_{35})\,]T_{32} = 0$ | 470 |
| $(b_{33})^{(6)}T_{32} - [(b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}(G_{35})\,]T_{33} = 0$ | 471 |
| $(b_{34})^{(6)}T_{33} - [(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}(G_{35})\,]T_{34} = 0$ | 472 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{36})^{(7)}G_{37} - \left[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37})\right]G_{36} = 0$ | 473 |
| $(a_{37})^{(7)}G_{36} - \left[(a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37})\right]G_{37} = 0$ | 474 |
| $(a_{38})^{(7)}G_{37} - \left[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37})\right]G_{38} = 0$ | 475 |
| $(b_{36})^{(7)}T_{37} - [(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}(G_{39})]T_{36} = 0$ | 476 |
| $(b_{37})^{(7)}T_{36} - [(b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}(G_{39})\,]T_{37} = 0$ | 477 |
| $(b_{38})^{(7)}T_{37} - [(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}(G_{39})\,]T_{38} = 0$ | 478 |
| $(a_{40})^{(8)}G_{41} - \left[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41})\right]G_{40} = 0$ | 479 |
| $(a_{41})^{(8)}G_{40} - \left[(a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41})\right]G_{41} = 0$ | 480 |
| $(a_{42})^{(8)}G_{41} - \left[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41})\right]G_{42} = 0$ | 481 |
| $(b_{40})^{(8)}T_{41} - [(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}(G_{43})\,]T_{40} = 0$ | 482 |
| $(b_{41})^{(8)}T_{40} - [(b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}(G_{43})\,]T_{41} = 0$ | 483 |

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| $(b_{42})^{(8)}T_{41} - [(b'_{42})^{(8)} - (b''_{42})^{(8)}(G_{43})\,]T_{42} = 0$ | 484 |
| $(a_{44})^{(9)}G_{45} - \left[(a'_{44})^{(9)} + (a''_{44})^{(9)}(T_{45})\right]G_{44} = 0$ | 484 A |
| $(a_{45})^{(9)}G_{44} - \left[(a'_{45})^{(9)} + (a''_{45})^{(9)}(T_{45})\right]G_{45} = 0$ | |
| $(a_{46})^{(9)}G_{45} - \left[(a'_{46})^{(9)} + (a''_{46})^{(9)}(T_{45})\right]G_{46} = 0$ | |
| $(b_{44})^{(9)}T_{45} - [(b'_{44})^{(9)} - (b''_{44})^{(9)}(G_{47})\,]T_{44} = 0$ | |
| $(b_{45})^{(9)}T_{44} - [(b'_{45})^{(9)} - (b''_{45})^{(9)}(G_{47})\,]T_{45} = 0$ | |
| $(b_{46})^{(9)}T_{45} - [(b'_{46})^{(9)} - (b''_{46})^{(9)}(G_{47})\,]T_{46} = 0$ | |
| **Proof:**<br>(a) Indeed the first two equations have a nontrivial solution $G_{13}, G_{14}$ if<br>$F(T) = (a'_{13})^{(1)}(a'_{14})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a'_{13})^{(1)}(a''_{14})^{(1)}(T_{14}) + (a'_{14})^{(1)}(a''_{13})^{(1)}(T_{14}) + (a''_{13})^{(1)}(T_{14})(a''_{14})^{(1)}(T_{14}) = 0$ | 485 |
| **Proof:**<br>(b) Indeed the first two equations have a nontrivial solution $G_{16}, G_{17}$ if<br>$\mathrm{F}(T_{19}) = (a'_{16})^{(2)}(a'_{17})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a'_{16})^{(2)}(a''_{17})^{(2)}(T_{17}) + (a'_{17})^{(2)}(a''_{16})^{(2)}(T_{17}) + (a''_{16})^{(2)}(T_{17})(a''_{17})^{(2)}(T_{17}) = 0$ | 486 |
| **Proof:**<br>(a) Indeed the first two equations have a nontrivial solution $G_{20}, G_{21}$ if<br>$F(T_{23}) = (a'_{20})^{(3)}(a'_{21})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a'_{20})^{(3)}(a''_{21})^{(3)}(T_{21}) + (a'_{21})^{(3)}(a''_{20})^{(3)}(T_{21}) + (a''_{20})^{(3)}(T_{21})(a''_{21})^{(3)}(T_{21}) = 0$ | 487 |
| **Proof:**<br>(a) Indeed the first two equations have a nontrivial solution $G_{24}, G_{25}$ if<br>$F(T_{27}) = (a'_{24})^{(4)}(a'_{25})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a'_{24})^{(4)}(a''_{25})^{(4)}(T_{25}) + (a'_{25})^{(4)}(a''_{24})^{(4)}(T_{25}) + (a''_{24})^{(4)}(T_{25})(a''_{25})^{(4)}(T_{25}) = 0$ | 488 |
| **Proof:**<br>(a) Indeed the first two equations have a nontrivial solution $G_{28}, G_{29}$ if<br>$F(T_{31}) = (a'_{28})^{(5)}(a'_{29})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a'_{28})^{(5)}(a''_{29})^{(5)}(T_{29}) + (a'_{29})^{(5)}(a''_{28})^{(5)}(T_{29}) +$ | 489 |

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| $(a_{28}^{''})^{(5)}(T_{29})(a_{29}^{''})^{(5)}(T_{29}) = 0$ | |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{32}, G_{33}$ if<br><br>$F(T_{35}) = (a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32}^{'})^{(6)}(a_{33}^{''})^{(6)}(T_{33}) + (a_{33}^{'})^{(6)}(a_{32}^{''})^{(6)}(T_{33}) + (a_{32}^{''})^{(6)}(T_{33})(a_{33}^{''})^{(6)}(T_{33}) = 0$ | 490 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{36}, G_{37}$ if<br><br>$F(T_{39}) = (a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36}^{'})^{(7)}(a_{37}^{''})^{(7)}(T_{37}) + (a_{37}^{'})^{(7)}(a_{36}^{''})^{(7)}(T_{37}) + (a_{36}^{''})^{(7)}(T_{37})(a_{37}^{''})^{(7)}(T_{37}) = 0$ | 491 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{40}, G_{41}$ if<br><br>$F(T_{43}) = (a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40}^{'})^{(8)}(a_{41}^{''})^{(8)}(T_{41}) + (a_{41}^{'})^{(8)}(a_{40}^{''})^{(8)}(T_{41}) + (a_{40}^{''})^{(8)}(T_{41})(a_{41}^{''})^{(8)}(T_{41}) = 0$ | 492 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{44}, G_{45}$ if<br><br>$F(T_{47}) = (a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44}^{'})^{(9)}(a_{45}^{''})^{(9)}(T_{45}) + (a_{45}^{'})^{(9)}(a_{44}^{''})^{(9)}(T_{45}) + (a_{44}^{''})^{(9)}(T_{45})(a_{45}^{''})^{(9)}(T_{45}) = 0$ | 492 A |
| **Definition and uniqueness of** $T_{14}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{14})$ being increasing, it follows that there exists a unique $T_{14}^{*}$ for which $f(T_{14}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{13} = \frac{(a_{13})^{(1)}G_{14}}{[(a_{13}^{'})^{(1)}+(a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15} = \frac{(a_{15})^{(1)}G_{14}}{[(a_{15}^{'})^{(1)}+(a_{15}^{''})^{(1)}(T_{14}^{*})]}$ | 493 |
| **Definition and uniqueness of** $T_{17}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(2)}(T_{17})$ being increasing, it follows that there exists a unique $T_{17}^{*}$ for which $f(T_{17}^{*}) = 0$. With this value , we obtain from the three first equations | 494 |
| $G_{16} = \frac{(a_{16})^{(2)}G_{17}}{[(a_{16}^{'})^{(2)}+(a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18} = \frac{(a_{18})^{(2)}G_{17}}{[(a_{18}^{'})^{(2)}+(a_{18}^{''})^{(2)}(T_{17}^{*})]}$ | 495 |
| **Definition and uniqueness of** $T_{21}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{21})$ being increasing, it follows that | 496 |

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| there exists a unique $T_{21}^{*}$ for which $f(T_{21}^{*})=0$. With this value , we obtain from the three first equations<br><br>$G_{20}=\frac{(a_{20})^{(3)}G_{21}}{[(a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21}^{*})]}$ , $G_{22}=\frac{(a_{22})^{(3)}G_{21}}{[(a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21}^{*})]}$ | |
| **Definition and uniqueness of** $T_{25}^{*}$ :-<br><br>After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(4)}(T_{25})$ being increasing, it follows that there exists a unique $T_{25}^{*}$ for which $f(T_{25}^{*})=0$. With this value , we obtain from the three first equations<br><br>$G_{24}=\frac{(a_{24})^{(4)}G_{25}}{[(a_{24}^{'})^{(4)}+(a_{24}^{''})^{(4)}(T_{25}^{*})]}$ , $G_{26}=\frac{(a_{26})^{(4)}G_{25}}{[(a_{26}^{'})^{(4)}+(a_{26}^{''})^{(4)}(T_{25}^{*})]}$ | 497 |
| **Definition and uniqueness of** $T_{29}^{*}$ :-<br><br>After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(5)}(T_{29})$ being increasing, it follows that there exists a unique $T_{29}^{*}$ for which $f(T_{29}^{*})=0$. With this value , we obtain from the three first equations<br><br>$G_{28}=\frac{(a_{28})^{(5)}G_{29}}{[(a_{28}^{'})^{(5)}+(a_{28}^{''})^{(5)}(T_{29}^{*})]}$ , $G_{30}=\frac{(a_{30})^{(5)}G_{29}}{[(a_{30}^{'})^{(5)}+(a_{30}^{''})^{(5)}(T_{29}^{*})]}$ | 498 |
| **Definition and uniqueness of** $T_{33}^{*}$ :-<br><br>After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(6)}(T_{33})$ being increasing, it follows that there exists a unique $T_{33}^{*}$ for which $f(T_{33}^{*})=0$. With this value , we obtain from the three first equations<br><br>$G_{32}=\frac{(a_{32})^{(6)}G_{33}}{[(a_{32}^{'})^{(6)}+(a_{32}^{''})^{(6)}(T_{33}^{*})]}$ , $G_{34}=\frac{(a_{34})^{(6)}G_{33}}{[(a_{34}^{'})^{(6)}+(a_{34}^{''})^{(6)}(T_{33}^{*})]}$ | 499 |
| **Definition and uniqueness of** $T_{37}^{*}$ :-<br><br>After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(7)}(T_{37})$ being increasing, it follows that there exists a unique $T_{37}^{*}$ for which $f(T_{37}^{*})=0$. With this value , we obtain from the three first equations<br><br>$G_{36}=\frac{(a_{36})^{(7)}G_{37}}{[(a_{36}^{'})^{(7)}+(a_{36}^{''})^{(7)}(T_{37}^{*})]}$ , $G_{38}=\frac{(a_{38})^{(7)}G_{37}}{[(a_{38}^{'})^{(7)}+(a_{38}^{''})^{(7)}(T_{37}^{*})]}$ | 500 |
| **Definition and uniqueness of** $T_{41}^{*}$ :-<br><br>After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(8)}(T_{41})$ being increasing, it follows that there exists a unique $T_{41}^{*}$ for which $f(T_{41}^{*})=0$. With this value , we obtain from the three first equations | 501 |

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| $G_{40} = \frac{(a_{40})^{(8)}G_{41}}{[(a_{40}^{'})^{(8)}+(a_{40}^{''})^{(8)}(T_{41}^{*})]}$ , $G_{42} = \frac{(a_{42})^{(8)}G_{41}}{[(a_{42}^{'})^{(8)}+(a_{42}^{''})^{(8)}(T_{41}^{*})]}$ | |
| **Definition and uniqueness of** $T_{45}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(9)}(T_{45})$ being increasing, it follows that there exists a unique $T_{45}^{*}$ for which $f(T_{45}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{44} = \frac{(a_{44})^{(9)}G_{45}}{[(a_{44}^{'})^{(9)}+(a_{44}^{''})^{(9)}(T_{45}^{*})]}$ , $G_{46} = \frac{(a_{46})^{(9)}G_{45}}{[(a_{46}^{'})^{(9)}+(a_{46}^{''})^{(9)}(T_{45}^{*})]}$ | 501 A |
| By the same argument, the equations admit solutions $G_{13}, G_{14}$ if<br>$\varphi(G) = (b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} -$<br>$\left[(b_{13}^{'})^{(1)}(b_{14}^{''})^{(1)}(G) + (b_{14}^{'})^{(1)}(b_{13}^{''})^{(1)}(G)\right]+(b_{13}^{''})^{(1)}(G)(b_{14}^{''})^{(1)}(G) = 0$<br><br>Where in $G(G_{13}, G_{14}, G_{15}), G_{13}, G_{15}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{14}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{14}^{*}$ such that $\varphi(G^{*}) = 0$ | 502 |
| By the same argument, the equations admit solutions $G_{16}, G_{17}$ if<br><br>$\varphi(G_{19}) = (b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} -$<br><br>$\left[(b_{16}^{'})^{(2)}(b_{17}^{''})^{(2)}(G_{19}) + (b_{17}^{'})^{(2)}(b_{16}^{''})^{(2)}(G_{19})\right]+(b_{16}^{''})^{(2)}(G_{19})(b_{17}^{''})^{(2)}(G_{19}) = 0$ | 503 |
| Where in $(G_{19})(G_{16}, G_{17}, G_{18}), G_{16}, G_{18}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{17}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{14}^{*}$ such that $\varphi((G_{19})^{*}) = 0$ | 504 |
| By the same argument, the equations admit solutions $G_{20}, G_{21}$ if<br><br>$\varphi(G_{23}) = (b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} -$<br><br>$\left[(b_{20}^{'})^{(3)}(b_{21}^{''})^{(3)}(G_{23}) + (b_{21}^{'})^{(3)}(b_{20}^{''})^{(3)}(G_{23})\right]+(b_{20}^{''})^{(3)}(G_{23})(b_{21}^{''})^{(3)}(G_{23}) = 0$<br><br>Where in $G_{23}(G_{20}, G_{21}, G_{22}), G_{20}, G_{22}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{21}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{21}^{*}$ such that $\varphi((G_{23})^{*}) = 0$ | 505 |
| By the same argument, the equations admit solutions $G_{24}, G_{25}$ if<br><br>$\varphi(G_{27}) = (b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} -$<br><br>$\left[(b_{24}^{'})^{(4)}(b_{25}^{''})^{(4)}(G_{27}) + (b_{25}^{'})^{(4)}(b_{24}^{''})^{(4)}(G_{27})\right]+(b_{24}^{''})^{(4)}(G_{27})(b_{25}^{''})^{(4)}(G_{27}) = 0$<br><br>Where in $(G_{27})(G_{24}, G_{25}, G_{26}), G_{24}, G_{26}$ must be replaced by their values from 96. It is easy to see that φ | 506 |

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| is a decreasing function in $G_{25}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{25}^{*}$ such that $\varphi((G_{27})^{*}) = 0$ | |
| By the same argument, the equations admit solutions $G_{28}, G_{29}$ if<br>$\varphi(G_{31}) = (b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} -$<br>$\left[(b_{28}^{'})^{(5)}(b_{29}^{''})^{(5)}(G_{31}) + (b_{29}^{'})^{(5)}(b_{28}^{''})^{(5)}(G_{31})\right] + (b_{28}^{''})^{(5)}(G_{31})(b_{29}^{''})^{(5)}(G_{31}) = 0$<br>Where in $(G_{31})(G_{28}, G_{29}, G_{30}), G_{28}, G_{30}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{29}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{29}^{*}$ such that $\varphi((G_{31})^{*}) = 0$ | 507 |
| By the same argument, the equations admit solutions $G_{32}, G_{33}$ if<br>$\varphi(G_{35}) = (b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} -$<br>$\left[(b_{32}^{'})^{(6)}(b_{33}^{''})^{(6)}(G_{35}) + (b_{33}^{'})^{(6)}(b_{32}^{''})^{(6)}(G_{35})\right] + (b_{32}^{''})^{(6)}(G_{35})(b_{33}^{''})^{(6)}(G_{35}) = 0$<br>Where in $(G_{35})(G_{32}, G_{33}, G_{34}), G_{32}, G_{34}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{33}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{33}^{*}$ such that $\varphi({G_{35}}^{*}) = 0$ | 508 |
| By the same argument, the equations admit solutions $G_{36}, G_{37}$ if<br>$\varphi(G_{39}) = (b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} -$<br>$\left[(b_{36}^{'})^{(7)}(b_{37}^{''})^{(7)}(G_{39}) + (b_{37}^{'})^{(7)}(b_{36}^{''})^{(7)}(G_{39})\right] + (b_{36}^{''})^{(7)}(G_{39})(b_{37}^{''})^{(7)}(G_{39}) = 0$<br>Where in $(G_{39})(G_{36}, G_{37}, G_{38}), G_{36}, G_{38}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{37}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{37}^{*}$ such that $\varphi({G_{39}}^{*}) = 0$ | 509 |
| By the same argument, the equations admit solutions $G_{40}, G_{41}$ if<br>$\varphi(G_{43}) = (b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} -$<br>$\left[(b_{40}^{'})^{(8)}(b_{41}^{''})^{(8)}(G_{43}) + (b_{41}^{'})^{(8)}(b_{40}^{''})^{(8)}(G_{43})\right] + (b_{40}^{''})^{(8)}(G_{43})(b_{41}^{''})^{(8)}(G_{43}) = 0$<br>Where in $(G_{43})(G_{40}, G_{41}, G_{42}), G_{40}, G_{42}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{41}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{41}^{*}$ such that $\varphi({G_{43}}^{*}) = 0$ | 510 |
| By the same argument, the equations 92,93 admit solutions $G_{44}, G_{45}$ if<br>$\varphi(G_{47}) = (b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} -$<br>$\left[(b_{44}^{'})^{(9)}(b_{45}^{''})^{(9)}(G_{47}) + (b_{45}^{'})^{(9)}(b_{44}^{''})^{(9)}(G_{47})\right] + (b_{44}^{''})^{(9)}(G_{47})(b_{45}^{''})^{(9)}(G_{47}) = 0$<br>Where in $(G_{47})(G_{44}, G_{45}, G_{46}), G_{44}, G_{46}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{45}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{45}^{*}$ such that $\varphi((G_{47})^{*}) = 0$ | |
| Finally we obtain the unique solution | 511 |

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| $G_{14}^{*}$ given by $\varphi(G^{*})=0$ , $T_{14}^{*}$ given by $f(T_{14}^{*})=0$ and<br>$G_{13}^{*}=\frac{(a_{13})^{(1)}G_{14}^{*}}{[(a_{13}^{'})^{(1)}+(a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15}^{*}=\frac{(a_{15})^{(1)}G_{14}^{*}}{[(a_{15}^{'})^{(1)}+(a_{15}^{''})^{(1)}(T_{14}^{*})]}$<br>$T_{13}^{*}=\frac{(b_{13})^{(1)}T_{14}^{*}}{[(b_{13}^{'})^{(1)}-(b_{13}^{''})^{(1)}(G^{*})]}$ , $T_{15}^{*}=\frac{(b_{15})^{(1)}T_{14}^{*}}{[(b_{15}^{'})^{(1)}-(b_{15}^{''})^{(1)}(G^{*})]}$<br>Obviously, these values represent an equilibrium solution | |
| Finally we obtain the unique solution | |
| $G_{17}^{*}$ given by $\varphi((G_{19})^{*})=0$ , $T_{17}^{*}$ given by $f(T_{17}^{*})=0$ and | 512 |
| $G_{16}^{*}=\frac{(a_{16})^{(2)}G_{17}^{*}}{[(a_{16}^{'})^{(2)}+(a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18}^{*}=\frac{(a_{18})^{(2)}G_{17}^{*}}{[(a_{18}^{'})^{(2)}+(a_{18}^{''})^{(2)}(T_{17}^{*})]}$ | 513 |
| $T_{16}^{*}=\frac{(b_{16})^{(2)}T_{17}^{*}}{[(b_{16}^{'})^{(2)}-(b_{16}^{''})^{(2)}((G_{19})^{*})]}$ , $T_{18}^{*}=\frac{(b_{18})^{(2)}T_{17}^{*}}{[(b_{18}^{'})^{(2)}-(b_{18}^{''})^{(2)}((G_{19})^{*})]}$ | 514 |
| Obviously, these values represent an equilibrium solution | |
| Finally we obtain the unique solution<br>$G_{21}^{*}$ given by $\varphi((G_{23})^{*})=0$ , $T_{21}^{*}$ given by $f(T_{21}^{*})=0$ and<br>$G_{20}^{*}=\frac{(a_{20})^{(3)}G_{21}^{*}}{[(a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21}^{*})]}$ , $G_{22}^{*}=\frac{(a_{22})^{(3)}G_{21}^{*}}{[(a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21}^{*})]}$<br>$T_{20}^{*}=\frac{(b_{20})^{(3)}T_{21}^{*}}{[(b_{20}^{'})^{(3)}-(b_{20}^{''})^{(3)}(G_{23}{}^{*})]}$ , $T_{22}^{*}=\frac{(b_{22})^{(3)}T_{21}^{*}}{[(b_{22}^{'})^{(3)}-(b_{22}^{''})^{(3)}(G_{23}{}^{*})]}$<br>Obviously, these values represent an equilibrium solution of global equations | 515 |
| Finally we obtain the unique solution<br>$G_{25}^{*}$ given by $\varphi(G_{27})=0$ , $T_{25}^{*}$ given by $f(T_{25}^{*})=0$ and<br>$G_{24}^{*}=\frac{(a_{24})^{(4)}G_{25}^{*}}{[(a_{24}^{'})^{(4)}+(a_{24}^{''})^{(4)}(T_{25}^{*})]}$ , $G_{26}^{*}=\frac{(a_{26})^{(4)}G_{25}^{*}}{[(a_{26}^{'})^{(4)}+(a_{26}^{''})^{(4)}(T_{25}^{*})]}$ | 516 |
| $T_{24}^{*}=\frac{(b_{24})^{(4)}T_{25}^{*}}{[(b_{24}^{'})^{(4)}-(b_{24}^{''})^{(4)}((G_{27})^{*})]}$ , $T_{26}^{*}=\frac{(b_{26})^{(4)}T_{25}^{*}}{[(b_{26}^{'})^{(4)}-(b_{26}^{''})^{(4)}((G_{27})^{*})]}$<br>Obviously, these values represent an equilibrium solution of global equations | 517 |
| Finally we obtain the unique solution<br>$G_{29}^{*}$ given by $\varphi((G_{31})^{*})=0$ , $T_{29}^{*}$ given by $f(T_{29}^{*})=0$ and<br>$G_{28}^{*}=\frac{(a_{28})^{(5)}G_{29}^{*}}{[(a_{28}^{'})^{(5)}+(a_{28}^{''})^{(5)}(T_{29}^{*})]}$ , $G_{30}^{*}=\frac{(a_{30})^{(5)}G_{29}^{*}}{[(a_{30}^{'})^{(5)}+(a_{30}^{''})^{(5)}(T_{29}^{*})]}$ | 518 |

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|---|---|
| $T_{28}^{*} = \frac{(b_{28})^{(5)}T_{29}^{*}}{[(b_{28}^{'})^{(5)}-(b_{28}^{''})^{(5)}((G_{31})^{*})]}$ , $T_{30}^{*} = \frac{(b_{30})^{(5)}T_{29}^{*}}{[(b_{30}^{'})^{(5)}-(b_{30}^{''})^{(5)}((G_{31})^{*})]}$<br>Obviously, these values represent an equilibrium solution of global equations | 519 |
| Finally we obtain the unique solution<br>$G_{33}^{*}$ given by $\varphi((G_{35})^{*}) = 0$ , $T_{33}^{*}$ given by $f(T_{33}^{*}) = 0$ and<br>$G_{32}^{*} = \frac{(a_{32})^{(6)}G_{33}^{*}}{[(a_{32}^{'})^{(6)}+(a_{32}^{''})^{(6)}(T_{33}^{*})]}$ , $G_{34}^{*} = \frac{(a_{34})^{(6)}G_{33}^{*}}{[(a_{34}^{'})^{(6)}+(a_{34}^{''})^{(6)}(T_{33}^{*})]}$ | 520 |
| $T_{32}^{*} = \frac{(b_{32})^{(6)}T_{33}^{*}}{[(b_{32}^{'})^{(6)}-(b_{32}^{''})^{(6)}((G_{35})^{*})]}$ , $T_{34}^{*} = \frac{(b_{34})^{(6)}T_{33}^{*}}{[(b_{34}^{'})^{(6)}-(b_{34}^{''})^{(6)}((G_{35})^{*})]}$<br>Obviously, these values represent an equilibrium solution of global equations | 521 |
| Finally we obtain the unique solution<br>$G_{37}^{*}$ given by $\varphi((G_{39})^{*}) = 0$ , $T_{37}^{*}$ given by $f(T_{37}^{*}) = 0$ and<br>$G_{36}^{*} = \frac{(a_{36})^{(7)}G_{37}^{*}}{[(a_{36}^{'})^{(7)}+(a_{36}^{''})^{(7)}(T_{37}^{*})]}$ , $G_{38}^{*} = \frac{(a_{38})^{(7)}G_{37}^{*}}{[(a_{38}^{'})^{(7)}+(a_{38}^{''})^{(7)}(T_{37}^{*})]}$<br>$T_{36}^{*} = \frac{(b_{36})^{(7)}T_{37}^{*}}{[(b_{36}^{'})^{(7)}-(b_{36}^{''})^{(7)}((G_{39})^{*})]}$ , $T_{38}^{*} = \frac{(b_{38})^{(7)}T_{37}^{*}}{[(b_{38}^{'})^{(7)}-(b_{38}^{''})^{(7)}((G_{39})^{*})]}$ | 522 |
| Finally we obtain the unique solution<br>$G_{41}^{*}$ given by $\varphi((G_{43})^{*}) = 0$ , $T_{41}^{*}$ given by $f(T_{41}^{*}) = 0$ and<br>$G_{40}^{*} = \frac{(a_{40})^{(8)}G_{41}^{*}}{[(a_{40}^{'})^{(8)}+(a_{40}^{''})^{(8)}(T_{41}^{*})]}$ , $G_{42}^{*} = \frac{(a_{42})^{(8)}G_{41}^{*}}{[(a_{42}^{'})^{(8)}+(a_{42}^{''})^{(8)}(T_{41}^{*})]}$<br>$T_{40}^{*} = \frac{(b_{40})^{(8)}T_{41}^{*}}{[(b_{40}^{'})^{(8)}-(b_{40}^{''})^{(8)}((G_{43})^{*})]}$ , $T_{42}^{*} = \frac{(b_{42})^{(8)}T_{41}^{*}}{[(b_{42}^{'})^{(8)}-(b_{42}^{''})^{(8)}((G_{43})^{*})]}$ | 523 |
| Finally we obtain the unique solution of 89 to 99<br>$G_{45}^{*}$ given by $\varphi((G_{47})^{*}) = 0$ , $T_{45}^{*}$ given by $f(T_{45}^{*}) = 0$ and<br>$G_{44}^{*} = \frac{(a_{44})^{(9)}G_{45}^{*}}{[(a_{44}^{'})^{(9)}+(a_{44}^{''})^{(9)}(T_{45}^{*})]}$ , $G_{46}^{*} = \frac{(a_{46})^{(9)}G_{45}^{*}}{[(a_{46}^{'})^{(9)}+(a_{46}^{''})^{(9)}(T_{45}^{*})]}$<br>$T_{44}^{*} = \frac{(b_{44})^{(9)}T_{45}^{*}}{[(b_{44}^{'})^{(9)}-(b_{44}^{''})^{(9)}((G_{47})^{*})]}$ , $T_{46}^{*} = \frac{(b_{46})^{(9)}T_{45}^{*}}{[(b_{46}^{'})^{(9)}-(b_{46}^{''})^{(9)}((G_{47})^{*})]}$ | 523 A |

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| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(1)}$ *and* $(b_i^{''})^{(1)}$ Belong to $C^{(1)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>**Proof:** Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial (a_{14}^{''})^{(1)}}{\partial T_{14}}(T_{14}^*) = (q_{14})^{(1)} \quad , \frac{\partial (b_i^{''})^{(1)}}{\partial G_j}(G^*) = s_{ij}$ | 524 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{13}}{dt} = -\left((a_{13}^{'})^{(1)} + (p_{13})^{(1)}\right)\mathbb{G}_{13} + (a_{13})^{(1)}\mathbb{G}_{14} - (q_{13})^{(1)}G_{13}^*\mathbb{T}_{14}$ | 525 |
| $\frac{d\mathbb{G}_{14}}{dt} = -\left((a_{14}^{'})^{(1)} + (p_{14})^{(1)}\right)\mathbb{G}_{14} + (a_{14})^{(1)}\mathbb{G}_{13} - (q_{14})^{(1)}G_{14}^*\mathbb{T}_{14}$ | 526 |
| $\frac{d\mathbb{G}_{15}}{dt} = -\left((a_{15}^{'})^{(1)} + (p_{15})^{(1)}\right)\mathbb{G}_{15} + (a_{15})^{(1)}\mathbb{G}_{14} - (q_{15})^{(1)}G_{15}^*\mathbb{T}_{14}$ | 527 |
| $\frac{d\mathbb{T}_{13}}{dt} = -\left((b_{13}^{'})^{(1)} - (r_{13})^{(1)}\right)\mathbb{T}_{13} + (b_{13})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(13)(j)}T_{13}^*\mathbb{G}_j\right)$ | 528 |
| $\frac{d\mathbb{T}_{14}}{dt} = -\left((b_{14}^{'})^{(1)} - (r_{14})^{(1)}\right)\mathbb{T}_{14} + (b_{14})^{(1)}\mathbb{T}_{13} + \sum_{j=13}^{15}\left(s_{(14)(j)}T_{14}^*\mathbb{G}_j\right)$ | 529 |
| $\frac{d\mathbb{T}_{15}}{dt} = -\left((b_{15}^{'})^{(1)} - (r_{15})^{(1)}\right)\mathbb{T}_{15} + (b_{15})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(15)(j)}T_{15}^*\mathbb{G}_j\right)$ | 530 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}$ Belong to $C^{(2)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable | 531 |
| Proof: Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- | |
| $G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$ | 532 |
| $\frac{\partial (a_{17}^{''})^{(2)}}{\partial T_{17}}(T_{17}^*) = (q_{17})^{(2)} \quad , \frac{\partial (b_i^{''})^{(2)}}{\partial G_j}((G_{19})^*) = s_{ij}$ | 533 |
| taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{16}}{dt} = -\left((a_{16}^{'})^{(2)} + (p_{16})^{(2)}\right)\mathbb{G}_{16} + (a_{16})^{(2)}\mathbb{G}_{17} - (q_{16})^{(2)}G_{16}^*\mathbb{T}_{17}$ | 534 |
| $\frac{d\mathbb{G}_{17}}{dt} = -\left((a_{17}^{'})^{(2)} + (p_{17})^{(2)}\right)\mathbb{G}_{17} + (a_{17})^{(2)}\mathbb{G}_{16} - (q_{17})^{(2)}G_{17}^*\mathbb{T}_{17}$ | 535 |

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| $\frac{d\mathbb{G}_{18}}{dt} = -\left((a_{18}^{'})^{(2)} + (p_{18})^{(2)}\right)\mathbb{G}_{18} + (a_{18})^{(2)}\mathbb{G}_{17} - (q_{18})^{(2)}G_{18}^{*}\mathbb{T}_{17}$ | 536 |
| $\frac{d\mathbb{T}_{16}}{dt} = -\left((b_{16}^{'})^{(2)} - (r_{16})^{(2)}\right)\mathbb{T}_{16} + (b_{16})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(16)(j)}T_{16}^{*}\mathbb{G}_{j}\right)$ | 537 |
| $\frac{d\mathbb{T}_{17}}{dt} = -\left((b_{17}^{'})^{(2)} - (r_{17})^{(2)}\right)\mathbb{T}_{17} + (b_{17})^{(2)}\mathbb{T}_{16} + \sum_{j=16}^{18}\left(s_{(17)(j)}T_{17}^{*}\mathbb{G}_{j}\right)$ | 538 |
| $\frac{d\mathbb{T}_{18}}{dt} = -\left((b_{18}^{'})^{(2)} - (r_{18})^{(2)}\right)\mathbb{T}_{18} + (b_{18})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(18)(j)}T_{18}^{*}\mathbb{G}_{j}\right)$ | 539 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(3)}$ *and* $(b_i^{''})^{(3)}$ Belong to $C^{(3)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^{*} + \mathbb{G}_i \quad , T_i = T_i^{*} + \mathbb{T}_i$<br><br>$\frac{\partial(a_{21}^{''})^{(3)}}{\partial T_{21}}(T_{21}^{*}) = (q_{21})^{(3)} \quad , \frac{\partial(b_i^{''})^{(3)}}{\partial G_j}((G_{23})^{*}) = s_{ij}$ | 540 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{20}}{dt} = -\left((a_{20}^{'})^{(3)} + (p_{20})^{(3)}\right)\mathbb{G}_{20} + (a_{20})^{(3)}\mathbb{G}_{21} - (q_{20})^{(3)}G_{20}^{*}\mathbb{T}_{21}$ | 541 |
| $\frac{d\mathbb{G}_{21}}{dt} = -\left((a_{21}^{'})^{(3)} + (p_{21})^{(3)}\right)\mathbb{G}_{21} + (a_{21})^{(3)}\mathbb{G}_{20} - (q_{21})^{(3)}G_{21}^{*}\mathbb{T}_{21}$ | 542 |
| $\frac{d\mathbb{G}_{22}}{dt} = -\left((a_{22}^{'})^{(3)} + (p_{22})^{(3)}\right)\mathbb{G}_{22} + (a_{22})^{(3)}\mathbb{G}_{21} - (q_{22})^{(3)}G_{22}^{*}\mathbb{T}_{21}$ | 543 |
| $\frac{d\mathbb{T}_{20}}{dt} = -\left((b_{20}^{'})^{(3)} - (r_{20})^{(3)}\right)\mathbb{T}_{20} + (b_{20})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(20)(j)}T_{20}^{*}\mathbb{G}_{j}\right)$ | 544 |
| $\frac{d\mathbb{T}_{21}}{dt} = -\left((b_{21}^{'})^{(3)} - (r_{21})^{(3)}\right)\mathbb{T}_{21} + (b_{21})^{(3)}\mathbb{T}_{20} + \sum_{j=20}^{22}\left(s_{(21)(j)}T_{21}^{*}\mathbb{G}_{j}\right)$ | 545 |
| $\frac{d\mathbb{T}_{22}}{dt} = -\left((b_{22}^{'})^{(3)} - (r_{22})^{(3)}\right)\mathbb{T}_{22} + (b_{22})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(22)(j)}T_{22}^{*}\mathbb{G}_{j}\right)$ | 546 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(4)}$ *and* $(b_i^{''})^{(4)}$ Belong to $C^{(4)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 547 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^{*} + \mathbb{G}_i \quad , T_i = T_i^{*} + \mathbb{T}_i$ | 548 |

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| $\frac{\partial (a_{25}^{''})^{(4)}}{\partial T_{25}}(T_{25}^{*}) = (q_{25})^{(4)}$ , $\frac{\partial (b_i^{''})^{(4)}}{\partial G_j}((G_{27})^{*}) = s_{ij}$ | |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{24}}{dt} = -\left((a_{24}^{'})^{(4)} + (p_{24})^{(4)}\right)\mathbb{G}_{24} + (a_{24})^{(4)}\mathbb{G}_{25} - (q_{24})^{(4)}G_{24}^{*}\mathbb{T}_{25}$ | 549 |
| $\frac{d\mathbb{G}_{25}}{dt} = -\left((a_{25}^{'})^{(4)} + (p_{25})^{(4)}\right)\mathbb{G}_{25} + (a_{25})^{(4)}\mathbb{G}_{24} - (q_{25})^{(4)}G_{25}^{*}\mathbb{T}_{25}$ | 550 |
| $\frac{d\mathbb{G}_{26}}{dt} = -\left((a_{26}^{'})^{(4)} + (p_{26})^{(4)}\right)\mathbb{G}_{26} + (a_{26})^{(4)}\mathbb{G}_{25} - (q_{26})^{(4)}G_{26}^{*}\mathbb{T}_{25}$ | 551 |
| $\frac{d\mathbb{T}_{24}}{dt} = -\left((b_{24}^{'})^{(4)} - (r_{24})^{(4)}\right)\mathbb{T}_{24} + (b_{24})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(24)(j)}T_{24}^{*}\mathbb{G}_j\right)$ | 552 |
| $\frac{d\mathbb{T}_{25}}{dt} = -\left((b_{25}^{'})^{(4)} - (r_{25})^{(4)}\right)\mathbb{T}_{25} + (b_{25})^{(4)}\mathbb{T}_{24} + \sum_{j=24}^{26}\left(s_{(25)(j)}T_{25}^{*}\mathbb{G}_j\right)$ | 553 |
| $\frac{d\mathbb{T}_{26}}{dt} = -\left((b_{26}^{'})^{(4)} - (r_{26})^{(4)}\right)\mathbb{T}_{26} + (b_{26})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(26)(j)}T_{26}^{*}\mathbb{G}_j\right)$ | 554 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 5:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(5)}$ *and* $(b_i^{''})^{(5)}$ Belong to $C^{(5)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 555 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^{*} + \mathbb{G}_i \quad , T_i = T_i^{*} + \mathbb{T}_i$<br><br>$\frac{\partial (a_{29}^{''})^{(5)}}{\partial T_{29}}(T_{29}^{*}) = (q_{29})^{(5)}$ , $\frac{\partial (b_i^{''})^{(5)}}{\partial G_j}((G_{31})^{*}) = s_{ij}$ | 556 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{28}}{dt} = -\left((a_{28}^{'})^{(5)} + (p_{28})^{(5)}\right)\mathbb{G}_{28} + (a_{28})^{(5)}\mathbb{G}_{29} - (q_{28})^{(5)}G_{28}^{*}\mathbb{T}_{29}$ | 557 |
| $\frac{d\mathbb{G}_{29}}{dt} = -\left((a_{29}^{'})^{(5)} + (p_{29})^{(5)}\right)\mathbb{G}_{29} + (a_{29})^{(5)}\mathbb{G}_{28} - (q_{29})^{(5)}G_{29}^{*}\mathbb{T}_{29}$ | 558 |
| $\frac{d\mathbb{G}_{30}}{dt} = -\left((a_{30}^{'})^{(5)} + (p_{30})^{(5)}\right)\mathbb{G}_{30} + (a_{30})^{(5)}\mathbb{G}_{29} - (q_{30})^{(5)}G_{30}^{*}\mathbb{T}_{29}$ | 559 |
| $\frac{d\mathbb{T}_{28}}{dt} = -\left((b_{28}^{'})^{(5)} - (r_{28})^{(5)}\right)\mathbb{T}_{28} + (b_{28})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(28)(j)}T_{28}^{*}\mathbb{G}_j\right)$ | 560 |
| $\frac{d\mathbb{T}_{29}}{dt} = -\left((b_{29}^{'})^{(5)} - (r_{29})^{(5)}\right)\mathbb{T}_{29} + (b_{29})^{(5)}\mathbb{T}_{28} + \sum_{j=28}^{30}\left(s_{(29)(j)}T_{29}^{*}\mathbb{G}_j\right)$ | 561 |
| $\frac{d\mathbb{T}_{30}}{dt} = -\left((b_{30}^{'})^{(5)} - (r_{30})^{(5)}\right)\mathbb{T}_{30} + (b_{30})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(30)(j)}T_{30}^{*}\mathbb{G}_j\right)$ | 562 |
| **ASYMPTOTIC STABILITY ANALYSIS** | 563 |

**Theorem 6:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(6)}$ $and$ $(b_i^{''})^{(6)}$ Belong to $C^{(6)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.

Proof: Denote

**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-

$$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$$

$$\frac{\partial (a_{33}^{''})^{(6)}}{\partial T_{33}}(T_{33}^*) = (q_{33})^{(6)} \quad , \frac{\partial (b_i^{''})^{(6)}}{\partial G_j}\left((G_{35})^*\right) = s_{ij}$$

Then taking into account equations and neglecting the terms of power 2, we obtain

$$\frac{d\mathbb{G}_{32}}{dt} = -\left((a_{32}^{'})^{(6)} + (p_{32})^{(6)}\right)\mathbb{G}_{32} + (a_{32})^{(6)}\mathbb{G}_{33} - (q_{32})^{(6)}G_{32}^*\mathbb{T}_{33} \quad (565)$$

$$\frac{d\mathbb{G}_{33}}{dt} = -\left((a_{33}^{'})^{(6)} + (p_{33})^{(6)}\right)\mathbb{G}_{33} + (a_{33})^{(6)}\mathbb{G}_{32} - (q_{33})^{(6)}G_{33}^*\mathbb{T}_{33} \quad (566)$$

$$\frac{d\mathbb{G}_{34}}{dt} = -\left((a_{34}^{'})^{(6)} + (p_{34})^{(6)}\right)\mathbb{G}_{34} + (a_{34})^{(6)}\mathbb{G}_{33} - (q_{34})^{(6)}G_{34}^*\mathbb{T}_{33} \quad (567)$$

$$\frac{d\mathbb{T}_{32}}{dt} = -\left((b_{32}^{'})^{(6)} - (r_{32})^{(6)}\right)\mathbb{T}_{32} + (b_{32})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(32)(j)}T_{32}^*\mathbb{G}_j\right) \quad (568)$$

$$\frac{d\mathbb{T}_{33}}{dt} = -\left((b_{33}^{'})^{(6)} - (r_{33})^{(6)}\right)\mathbb{T}_{33} + (b_{33})^{(6)}\mathbb{T}_{32} + \sum_{j=32}^{34}\left(s_{(33)(j)}T_{33}^*\mathbb{G}_j\right) \quad (569)$$

$$\frac{d\mathbb{T}_{34}}{dt} = -\left((b_{34}^{'})^{(6)} - (r_{34})^{(6)}\right)\mathbb{T}_{34} + (b_{34})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(34)(j)}T_{34}^*\mathbb{G}_j\right) \quad (570)$$

**ASYMPTOTIC STABILITY ANALYSIS** (571)

**Theorem 7:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(7)}$ $and$ $(b_i^{''})^{(7)}$ Belong to $C^{(7)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.

Proof: Denote

**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- (572)

$$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$$

$$\frac{\partial (a_{37}^{''})^{(7)}}{\partial T_{37}}(T_{37}^*) = (q_{37})^{(7)} \quad , \frac{\partial (b_i^{''})^{(7)}}{\partial G_j}\left((G_{39})^{**}\right) = s_{ij}$$

Then taking into account equations and neglecting the terms of power 2, we obtain from

$$\frac{d\mathbb{G}_{36}}{dt} = -\left((a_{36}^{'})^{(7)} + (p_{36})^{(7)}\right)\mathbb{G}_{36} + (a_{36})^{(7)}\mathbb{G}_{37} - (q_{36})^{(7)}G_{36}^*\mathbb{T}_{37} \quad (573)$$

$$\frac{d\mathbb{G}_{37}}{dt} = -\left((a_{37}^{'})^{(7)} + (p_{37})^{(7)}\right)\mathbb{G}_{37} + (a_{37})^{(7)}\mathbb{G}_{36} - (q_{37})^{(7)}G_{37}^*\mathbb{T}_{37} \quad (574)$$

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| $\frac{d\mathbb{G}_{38}}{dt} = -\left((a'_{38})^{(7)} + (p_{38})^{(7)}\right)\mathbb{G}_{38} + (a_{38})^{(7)}\mathbb{G}_{37} - (q_{38})^{(7)}G^*_{38}\mathbb{T}_{37}$ | 575 |
| $\frac{d\mathbb{T}_{36}}{dt} = -\left((b'_{36})^{(7)} - (r_{36})^{(7)}\right)\mathbb{T}_{36} + (b_{36})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(36)(j)}T^*_{36}\mathbb{G}_j\right)$ | 576 |
| $\frac{d\mathbb{T}_{37}}{dt} = -\left((b'_{37})^{(7)} - (r_{37})^{(7)}\right)\mathbb{T}_{37} + (b_{37})^{(7)}\mathbb{T}_{36} + \sum_{j=36}^{38}\left(s_{(37)(j)}T^*_{37}\mathbb{G}_j\right)$ | 578 |
| $\frac{d\mathbb{T}_{38}}{dt} = -\left((b'_{38})^{(7)} - (r_{38})^{(7)}\right)\mathbb{T}_{38} + (b_{38})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(38)(j)}T^*_{38}\mathbb{G}_j\right)$ | 579 |
| Obviously, these values represent an equilibrium solution<br><br>**ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 8:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(8)}$ *and* $(b_i^{''})^{(8)}$ Belong to $C^{(8)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial(a''_{41})^{(8)}}{\partial T_{41}}(T^*_{41}) = (q_{41})^{(8)}$ , $\frac{\partial(b_i^{''})^{(8)}}{\partial G_j}((G_{43})^*) = s_{ij}$ | 580 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{40}}{dt} = -\left((a'_{40})^{(8)} + (p_{40})^{(8)}\right)\mathbb{G}_{40} + (a_{40})^{(8)}\mathbb{G}_{41} - (q_{40})^{(8)}G^*_{40}\mathbb{T}_{41}$ | 581 |
| $\frac{d\mathbb{G}_{41}}{dt} = -\left((a'_{41})^{(8)} + (p_{41})^{(8)}\right)\mathbb{G}_{41} + (a_{41})^{(8)}\mathbb{G}_{40} - (q_{41})^{(8)}G^*_{41}\mathbb{T}_{41}$ | 582 |
| $\frac{d\mathbb{G}_{42}}{dt} = -\left((a'_{42})^{(8)} + (p_{42})^{(8)}\right)\mathbb{G}_{42} + (a_{42})^{(8)}\mathbb{G}_{41} - (q_{42})^{(8)}G^*_{42}\mathbb{T}_{41}$ | 583 |
| $\frac{d\mathbb{T}_{40}}{dt} = -\left((b'_{40})^{(8)} - (r_{40})^{(8)}\right)\mathbb{T}_{40} + (b_{40})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(40)(j)}T^*_{40}\mathbb{G}_j\right)$ | 584 |
| $\frac{d\mathbb{T}_{41}}{dt} = -\left((b'_{41})^{(8)} - (r_{41})^{(8)}\right)\mathbb{T}_{41} + (b_{41})^{(8)}\mathbb{T}_{40} + \sum_{j=40}^{42}\left(s_{(41)(j)}T^*_{41}\mathbb{G}_j\right)$ | 585 |
| $\frac{d\mathbb{T}_{42}}{dt} = -\left((b'_{42})^{(8)} - (r_{42})^{(8)}\right)\mathbb{T}_{42} + (b_{42})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(42)(j)}T^*_{42}\mathbb{G}_j\right)$ | 586 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 9:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(9)}$ *and* $(b_i^{''})^{(9)}$ Belong to $C^{(9)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br>**Proof: Denote** | 586 A |

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| | |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br>$\frac{\partial (a_{45}^{''})^{(9)}}{\partial T_{45}}(T_{45}^*) = (q_{45})^{(9)} \ , \ \frac{\partial (b_i^{''})^{(9)}}{\partial G_j}((G_{47})^*) = s_{ij}$ | |
| Then taking into account equations 89 to 99 and neglecting the terms of power 2, we obtain from 99 to 44 | |
| $\frac{d\mathbb{G}_{44}}{dt} = -\left((a_{44}^{'})^{(9)} + (p_{44})^{(9)}\right)\mathbb{G}_{44} + (a_{44})^{(9)}\mathbb{G}_{45} - (q_{44})^{(9)}G_{44}^*\mathbb{T}_{45}$ | 586 B |
| $\frac{d\mathbb{G}_{45}}{dt} = -\left((a_{45}^{'})^{(9)} + (p_{45})^{(9)}\right)\mathbb{G}_{45} + (a_{45})^{(9)}\mathbb{G}_{44} - (q_{45})^{(9)}G_{45}^*\mathbb{T}_{45}$ | 586 C |
| $\frac{d\mathbb{G}_{46}}{dt} = -\left((a_{46}^{'})^{(9)} + (p_{46})^{(9)}\right)\mathbb{G}_{46} + (a_{46})^{(9)}\mathbb{G}_{45} - (q_{46})^{(9)}G_{46}^*\mathbb{T}_{45}$ | 586 D |
| $\frac{d\mathbb{T}_{44}}{dt} = -\left((b_{44}^{'})^{(9)} - (r_{44})^{(9)}\right)\mathbb{T}_{44} + (b_{44})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(44)(j)}T_{44}^*\mathbb{G}_j\right)$ | 586 E |
| $\frac{d\mathbb{T}_{45}}{dt} = -\left((b_{45}^{'})^{(9)} - (r_{45})^{(9)}\right)\mathbb{T}_{45} + (b_{45})^{(9)}\mathbb{T}_{44} + \sum_{j=44}^{46}\left(s_{(45)(j)}T_{45}^*\mathbb{G}_j\right)$ | 586 F |
| $\frac{d\mathbb{T}_{46}}{dt} = -\left((b_{46}^{'})^{(9)} - (r_{46})^{(9)}\right)\mathbb{T}_{46} + (b_{46})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(46)(j)}T_{46}^*\mathbb{G}_j\right)$ | 586 G |
| **The characteristic equation of this system is** | 587 |
| $\left((\lambda)^{(1)} + (b_{15}^{'})^{(1)} - (r_{15})^{(1)}\right)\{\left((\lambda)^{(1)} + (a_{15}^{'})^{(1)} + (p_{15})^{(1)}\right)$<br>$\left[\left((\lambda)^{(1)} + (a_{13}^{'})^{(1)} + (p_{13})^{(1)}\right)(q_{14})^{(1)}G_{14}^* + (a_{14})^{(1)}(q_{13})^{(1)}G_{13}^*\right)\Big]$<br>$\left((\lambda)^{(1)} + (b_{13}^{'})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(14)}T_{14}^* + (b_{14})^{(1)}s_{(13),(14)}T_{14}^*\Big)$<br>$+\left(\left((\lambda)^{(1)} + (a_{14}^{'})^{(1)} + (p_{14})^{(1)}\right)(q_{13})^{(1)}G_{13}^* + (a_{13})^{(1)}(q_{14})^{(1)}G_{14}^*\right)$<br>$\left((\lambda)^{(1)} + (b_{13}^{'})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(13)}T_{14}^* + (b_{14})^{(1)}s_{(13),(13)}T_{13}^*\Big)$<br>$\left((\lambda)^{(1)}\right)^2 + \left((a_{13}^{'})^{(1)} + (a_{14}^{'})^{(1)} + (p_{13})^{(1)} + (p_{14})^{(1)}\right)(\lambda)^{(1)}\Big)$<br>$\left((\lambda)^{(1)}\right)^2 + \left((b_{13}^{'})^{(1)} + (b_{14}^{'})^{(1)} - (r_{13})^{(1)} + (r_{14})^{(1)}\right)(\lambda)^{(1)}\Big)$<br>$+\left(\left((\lambda)^{(1)}\right)^2 + \left((a_{13}^{'})^{(1)} + (a_{14}^{'})^{(1)} + (p_{13})^{(1)} + (p_{14})^{(1)}\right)(\lambda)^{(1)}\right)(q_{15})^{(1)}G_{15}$<br>$+\left((\lambda)^{(1)} + (a_{13}^{'})^{(1)} + (p_{13})^{(1)}\right)\left((a_{15})^{(1)}(q_{14})^{(1)}G_{14}^* + (a_{14})^{(1)}(a_{15})^{(1)}(q_{13})^{(1)}G_{13}^*\right)$<br>$\left((\lambda)^{(1)} + (b_{13}^{'})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(15)}T_{14}^* + (b_{14})^{(1)}s_{(13),(15)}T_{13}^*\Big)\} = 0$ | |

$$+$$

$$((\lambda)^{(2)} + (b'_{18})^{(2)} - (r_{18})^{(2)})\{((\lambda)^{(2)} + (a'_{18})^{(2)} + (p_{18})^{(2)})$$

$$\left[\left(((\lambda)^{(2)} + (a'_{16})^{(2)} + (p_{16})^{(2)})(q_{17})^{(2)}\mathrm{G}^{*}_{17} + (a_{17})^{(2)}(q_{16})^{(2)}\mathrm{G}^{*}_{16}\right)\right]$$

$$\left(((\lambda)^{(2)} + (b'_{16})^{(2)} - (r_{16})^{(2)})s_{(17),(17)}\mathrm{T}^{*}_{17} + (b_{17})^{(2)}s_{(16),(17)}\mathrm{T}^{*}_{17}\right)$$

$$+\left(((\lambda)^{(2)} + (a'_{17})^{(2)} + (p_{17})^{(2)})(q_{16})^{(2)}\mathrm{G}^{*}_{16} + (a_{16})^{(2)}(q_{17})^{(2)}\mathrm{G}^{*}_{17}\right)$$

$$\left(((\lambda)^{(2)} + (b'_{16})^{(2)} - (r_{16})^{(2)})s_{(17),(16)}\mathrm{T}^{*}_{17} + (b_{17})^{(2)}s_{(16),(16)}\mathrm{T}^{*}_{16}\right)$$

$$\left((\lambda)^{(2)}\right)^{2} + \left((a'_{16})^{(2)} + (a'_{17})^{(2)} + (p_{16})^{(2)} + (p_{17})^{(2)}\right)(\lambda)^{(2)}\Big)$$

$$\left((\lambda)^{(2)}\right)^{2} + \left((b'_{16})^{(2)} + (b'_{17})^{(2)} - (r_{16})^{(2)} + (r_{17})^{(2)}\right)(\lambda)^{(2)}\Big)$$

$$+\left(\left((\lambda)^{(2)}\right)^{2} + \left((a'_{16})^{(2)} + (a'_{17})^{(2)} + (p_{16})^{(2)} + (p_{17})^{(2)}\right)(\lambda)^{(2)}\right)(q_{18})^{(2)}\mathrm{G}_{18}$$

$$+\left((\lambda)^{(2)} + (a'_{16})^{(2)} + (p_{16})^{(2)}\right)\left((a_{18})^{(2)}(q_{17})^{(2)}\mathrm{G}^{*}_{17} + (a_{17})^{(2)}(a_{18})^{(2)}(q_{16})^{(2)}\mathrm{G}^{*}_{16}\right)$$

$$\left(((\lambda)^{(2)} + (b'_{16})^{(2)} - (r_{16})^{(2)})s_{(17),(18)}\mathrm{T}^{*}_{17} + (b_{17})^{(2)}s_{(16),(18)}\mathrm{T}^{*}_{16}\right)\} = 0$$

$$+$$

$$((\lambda)^{(3)} + (b'_{22})^{(3)} - (r_{22})^{(3)})\{((\lambda)^{(3)} + (a'_{22})^{(3)} + (p_{22})^{(3)})$$

$$\left[\left(((\lambda)^{(3)} + (a'_{20})^{(3)} + (p_{20})^{(3)})(q_{21})^{(3)}G^{*}_{21} + (a_{21})^{(3)}(q_{20})^{(3)}G^{*}_{20}\right)\right]$$

$$\left(((\lambda)^{(3)} + (b'_{20})^{(3)} - (r_{20})^{(3)})s_{(21),(21)}T^{*}_{21} + (b_{21})^{(3)}s_{(20),(21)}T^{*}_{21}\right)$$

$$+\left(((\lambda)^{(3)} + (a'_{21})^{(3)} + (p_{21})^{(3)})(q_{20})^{(3)}G^{*}_{20} + (a_{20})^{(3)}(q_{21})^{(1)}G^{*}_{21}\right)$$

$$\left(((\lambda)^{(3)} + (b'_{20})^{(3)} - (r_{20})^{(3)})s_{(21),(20)}T^{*}_{21} + (b_{21})^{(3)}s_{(20),(20)}T^{*}_{20}\right)$$

$$\left((\lambda)^{(3)}\right)^{2} + \left((a'_{20})^{(3)} + (a'_{21})^{(3)} + (p_{20})^{(3)} + (p_{21})^{(3)}\right)(\lambda)^{(3)}\Big)$$

$$\left((\lambda)^{(3)}\right)^{2} + \left((b'_{20})^{(3)} + (b'_{21})^{(3)} - (r_{20})^{(3)} + (r_{21})^{(3)}\right)(\lambda)^{(3)}\Big)$$

$$+\left(\left((\lambda)^{(3)}\right)^{2} + \left((a'_{20})^{(3)} + (a'_{21})^{(3)} + (p_{20})^{(3)} + (p_{21})^{(3)}\right)(\lambda)^{(3)}\right)(q_{22})^{(3)}G_{22}$$

$$+\left((\lambda)^{(3)} + (a'_{20})^{(3)} + (p_{20})^{(3)}\right)\left((a_{22})^{(3)}(q_{21})^{(3)}G^{*}_{21} + (a_{21})^{(3)}(a_{22})^{(3)}(q_{20})^{(3)}G^{*}_{20}\right)$$

$$\left(((\lambda)^{(3)} + (b'_{20})^{(3)} - (r_{20})^{(3)})s_{(21),(22)}T^{*}_{21} + (b_{21})^{(3)}s_{(20),(22)}T^{*}_{20}\right)\} = 0$$

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| $+$ | |
| $\left((\lambda)^{(4)}+(b_{26}^{'})^{(4)}-(r_{26})^{(4)}\right)\{\left((\lambda)^{(4)}+(a_{26}^{'})^{(4)}+(p_{26})^{(4)}\right)$ $\left[\left(\left((\lambda)^{(4)}+(a_{24}^{'})^{(4)}+(p_{24})^{(4)}\right)(q_{25})^{(4)}G_{25}^{*}+(a_{25})^{(4)}(q_{24})^{(4)}G_{24}^{*}\right)\right]$ $\left(\left((\lambda)^{(4)}+(b_{24}^{'})^{(4)}-(r_{24})^{(4)}\right)s_{(25),(25)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(25)}T_{25}^{*}\right)$ $+\left(\left((\lambda)^{(4)}+(a_{25}^{'})^{(4)}+(p_{25})^{(4)}\right)(q_{24})^{(4)}G_{24}^{*}+(a_{24})^{(4)}(q_{25})^{(4)}G_{25}^{*}\right)$ $\left(\left((\lambda)^{(4)}+(b_{24}^{'})^{(4)}-(r_{24})^{(4)}\right)s_{(25),(24)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(24)}T_{24}^{*}\right)$ $\left(\left((\lambda)^{(4)}\right)^{2}+\left((a_{24}^{'})^{(4)}+(a_{25}^{'})^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)$ $\left(\left((\lambda)^{(4)}\right)^{2}+\left((b_{24}^{'})^{(4)}+(b_{25}^{'})^{(4)}-(r_{24})^{(4)}+(r_{25})^{(4)}\right)(\lambda)^{(4)}\right)$ $+\left(\left((\lambda)^{(4)}\right)^{2}+\left((a_{24}^{'})^{(4)}+(a_{25}^{'})^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)(q_{26})^{(4)}G_{26}$ $+\left((\lambda)^{(4)}+(a_{24}^{'})^{(4)}+(p_{24})^{(4)}\right)\left((a_{26})^{(4)}(q_{25})^{(4)}G_{25}^{*}+(a_{25})^{(4)}(a_{26})^{(4)}(q_{24})^{(4)}G_{24}^{*}\right)$ $\left(\left((\lambda)^{(4)}+(b_{24}^{'})^{(4)}-(r_{24})^{(4)}\right)s_{(25),(26)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(26)}T_{24}^{*}\right)\}=0$ | |
| $+$ | |
| $\left((\lambda)^{(5)}+(b_{30}^{'})^{(5)}-(r_{30})^{(5)}\right)\{\left((\lambda)^{(5)}+(a_{30}^{'})^{(5)}+(p_{30})^{(5)}\right)$ $\left[\left(\left((\lambda)^{(5)}+(a_{28}^{'})^{(5)}+(p_{28})^{(5)}\right)(q_{29})^{(5)}G_{29}^{*}+(a_{29})^{(5)}(q_{28})^{(5)}G_{28}^{*}\right)\right]$ $\left(\left((\lambda)^{(5)}+(b_{28}^{'})^{(5)}-(r_{28})^{(5)}\right)s_{(29),(29)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(29)}T_{29}^{*}\right)$ $+\left(\left((\lambda)^{(5)}+(a_{29}^{'})^{(5)}+(p_{29})^{(5)}\right)(q_{28})^{(5)}G_{28}^{*}+(a_{28})^{(5)}(q_{29})^{(5)}G_{29}^{*}\right)$ $\left(\left((\lambda)^{(5)}+(b_{28}^{'})^{(5)}-(r_{28})^{(5)}\right)s_{(29),(28)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(28)}T_{28}^{*}\right)$ $\left(\left((\lambda)^{(5)}\right)^{2}+\left((a_{28}^{'})^{(5)}+(a_{29}^{'})^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)$ $\left(\left((\lambda)^{(5)}\right)^{2}+\left((b_{28}^{'})^{(5)}+(b_{29}^{'})^{(5)}-(r_{28})^{(5)}+(r_{29})^{(5)}\right)(\lambda)^{(5)}\right)$ $+\left(\left((\lambda)^{(5)}\right)^{2}+\left((a_{28}^{'})^{(5)}+(a_{29}^{'})^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)(q_{30})^{(5)}G_{30}$ $+\left((\lambda)^{(5)}+(a_{28}^{'})^{(5)}+(p_{28})^{(5)}\right)\left((a_{30})^{(5)}(q_{29})^{(5)}G_{29}^{*}+(a_{29})^{(5)}(a_{30})^{(5)}(q_{28})^{(5)}G_{28}^{*}\right)$ $\left(\left((\lambda)^{(5)}+(b_{28}^{'})^{(5)}-(r_{28})^{(5)}\right)s_{(29),(30)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(30)}T_{28}^{*}\right)\}=0$ | |

$$+$$

$$\left((\lambda)^{(6)}+(b_{34}^{'})^{(6)}-(r_{34})^{(6)}\right)\{\left((\lambda)^{(6)}+(a_{34}^{'})^{(6)}+(p_{34})^{(6)}\right)$$

$$\left[\left(\left((\lambda)^{(6)}+(a_{32}^{'})^{(6)}+(p_{32})^{(6)}\right)(q_{33})^{(6)}G_{33}^{*}+(a_{33})^{(6)}(q_{32})^{(6)}G_{32}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(6)}+(b_{32}^{'})^{(6)}-(r_{32})^{(6)}\right)s_{(33),(33)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(33)}T_{33}^{*}\right)$$

$$+\left(\left((\lambda)^{(6)}+(a_{33}^{'})^{(6)}+(p_{33})^{(6)}\right)(q_{32})^{(6)}G_{32}^{*}+(a_{32})^{(6)}(q_{33})^{(6)}G_{33}^{*}\right)$$

$$\left(\left((\lambda)^{(6)}+(b_{32}^{'})^{(6)}-(r_{32})^{(6)}\right)s_{(33),(32)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(32)}T_{32}^{*}\right)$$

$$\left(\left((\lambda)^{(6)}\right)^{2}+\left((a_{32}^{'})^{(6)}+(a_{33}^{'})^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\right)$$

$$\left(\left((\lambda)^{(6)}\right)^{2}+\left((b_{32}^{'})^{(6)}+(b_{33}^{'})^{(6)}-(r_{32})^{(6)}+(r_{33})^{(6)}\right)(\lambda)^{(6)}\right)$$

$$+\left(\left((\lambda)^{(6)}\right)^{2}+\left((a_{32}^{'})^{(6)}+(a_{33}^{'})^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\right)(q_{34})^{(6)}G_{34}$$

$$+\left((\lambda)^{(6)}+(a_{32}^{'})^{(6)}+(p_{32})^{(6)}\right)\left((a_{34})^{(6)}(q_{33})^{(6)}G_{33}^{*}+(a_{33})^{(6)}(a_{34})^{(6)}(q_{32})^{(6)}G_{32}^{*}\right)$$

$$\left(\left((\lambda)^{(6)}+(b_{32}^{'})^{(6)}-(r_{32})^{(6)}\right)s_{(33),(34)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(34)}T_{32}^{*}\right)\}=0$$

$$+$$

$$\left((\lambda)^{(7)}+(b_{38}^{'})^{(7)}-(r_{38})^{(7)}\right)\{\left((\lambda)^{(7)}+(a_{38}^{'})^{(7)}+(p_{38})^{(7)}\right)$$

$$\left[\left(\left((\lambda)^{(7)}+(a_{36}^{'})^{(7)}+(p_{36})^{(7)}\right)(q_{37})^{(7)}G_{37}^{*}+(a_{37})^{(7)}(q_{36})^{(7)}G_{36}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(7)}+(b_{36}^{'})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(37)}T_{37}^{*}+(b_{37})^{(7)}s_{(36),(37)}T_{37}^{*}\right)$$

$$+\left(\left((\lambda)^{(7)}+(a_{37}^{'})^{(7)}+(p_{37})^{(7)}\right)(q_{36})^{(7)}G_{36}^{*}+(a_{36})^{(7)}(q_{37})^{(7)}G_{37}^{*}\right)$$

$$\left(\left((\lambda)^{(7)}+(b_{36}^{'})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(36)}T_{37}^{*}+(b_{37})^{(7)}s_{(36),(36)}T_{36}^{*}\right)$$

$$\left(\left((\lambda)^{(7)}\right)^{2}+\left((a_{36}^{'})^{(7)}+(a_{37}^{'})^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\right)$$

$$\left(\left((\lambda)^{(7)}\right)^{2}+\left((b_{36}^{'})^{(7)}+(b_{37}^{'})^{(7)}-(r_{36})^{(7)}+(r_{37})^{(7)}\right)(\lambda)^{(7)}\right)$$

$$+\left(\left((\lambda)^{(7)}\right)^{2}+\left((a_{36}^{'})^{(7)}+(a_{37}^{'})^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\right)(q_{38})^{(7)}G_{38}$$

$$+\left((\lambda)^{(7)}+(a_{36}^{'})^{(7)}+(p_{36})^{(7)}\right)\left((a_{38})^{(7)}(q_{37})^{(7)}G_{37}^{*}+(a_{37})^{(7)}(a_{38})^{(7)}(q_{36})^{(7)}G_{36}^{*}\right)$$

$$\left(\left((\lambda)^{(7)}+(b_{36}^{'})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(38)}T_{37}^{*}+(b_{37})^{(7)}s_{(36),(38)}T_{36}^{*}\right)\}=0$$

$$+$$

$$\left((\lambda)^{(8)}+(b'_{42})^{(8)}-(r_{42})^{(8)}\right)\{\left((\lambda)^{(8)}+(a'_{42})^{(8)}+(p_{42})^{(8)}\right)$$

$$\left[\left(\left((\lambda)^{(8)}+(a'_{40})^{(8)}+(p_{40})^{(8)}\right)(q_{41})^{(8)}G^{*}_{41}+(a_{41})^{(8)}(q_{40})^{(8)}G^{*}_{40}\right)\right]$$

$$\left(\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(41)}T^{*}_{41}+(b_{41})^{(8)}s_{(40),(41)}T^{*}_{41}\right)$$

$$+\left(\left((\lambda)^{(8)}+(a'_{41})^{(8)}+(p_{41})^{(8)}\right)(q_{40})^{(8)}G^{*}_{40}+(a_{40})^{(8)}(q_{41})^{(8)}G^{*}_{41}\right)$$

$$\left(\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(40)}T^{*}_{41}+(b_{41})^{(8)}s_{(40),(40)}T^{*}_{40}\right)$$

$$\left(\left((\lambda)^{(8)}\right)^{2}+\left((a'_{40})^{(8)}+(a'_{41})^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\right)$$

$$\left(\left((\lambda)^{(8)}\right)^{2}+\left((b'_{40})^{(8)}+(b'_{41})^{(8)}-(r_{40})^{(8)}+(r_{41})^{(8)}\right)(\lambda)^{(8)}\right)$$

$$+\left(\left((\lambda)^{(8)}\right)^{2}+\left((a'_{40})^{(8)}+(a'_{41})^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\right)(q_{42})^{(8)}G_{42}$$

$$+\left((\lambda)^{(8)}+(a'_{40})^{(8)}+(p_{40})^{(8)}\right)\left((a_{42})^{(8)}(q_{41})^{(8)}G^{*}_{41}+(a_{41})^{(8)}(a_{42})^{(8)}(q_{40})^{(8)}G^{*}_{40}\right)$$

$$\left(\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(42)}T^{*}_{41}+(b_{41})^{(8)}s_{(40),(42)}T^{*}_{40}\right)\}=0$$

$$+$$

$$\left((\lambda)^{(9)}+(b'_{46})^{(9)}-(r_{46})^{(9)}\right)\{\left((\lambda)^{(9)}+(a'_{46})^{(9)}+(p_{46})^{(9)}\right)$$

$$\left[\left(\left((\lambda)^{(9)}+(a'_{44})^{(9)}+(p_{44})^{(9)}\right)(q_{45})^{(9)}G^{*}_{45}+(a_{45})^{(9)}(q_{44})^{(9)}G^{*}_{44}\right)\right]$$

$$\left(\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(45)}T^{*}_{45}+(b_{45})^{(9)}s_{(44),(45)}T^{*}_{45}\right)$$

$$+\left(\left((\lambda)^{(9)}+(a'_{45})^{(9)}+(p_{45})^{(9)}\right)(q_{44})^{(9)}G^{*}_{44}+(a_{44})^{(9)}(q_{45})^{(9)}G^{*}_{45}\right)$$

$$\left(\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(44)}T^{*}_{45}+(b_{45})^{(9)}s_{(44),(44)}T^{*}_{44}\right)$$

$$\left(\left((\lambda)^{(9)}\right)^{2}+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\right)$$

$$\left(\left((\lambda)^{(9)}\right)^{2}+\left((b'_{44})^{(9)}+(b'_{45})^{(9)}-(r_{44})^{(9)}+(r_{45})^{(9)}\right)(\lambda)^{(9)}\right)$$

$$+\left(\left((\lambda)^{(9)}\right)^{2}+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\right)(q_{46})^{(9)}G_{46}$$

$$+\left((\lambda)^{(9)}+(a'_{44})^{(9)}+(p_{44})^{(9)}\right)\left((a_{46})^{(9)}(q_{45})^{(9)}G^{*}_{45}+(a_{45})^{(9)}(a_{46})^{(9)}(q_{44})^{(9)}G^{*}_{44}\right)$$

$\left(\left((\lambda)^{(9)} + (b'_{44})^{(9)} - (r_{44})^{(9)}\right)s_{(45),(46)}T^*_{45} + (b_{45})^{(9)}s_{(44),(46)}T^*_{44}\right)\} = 0$

**And as one sees, all the coefficients are positive. It follows that all the roots have negative real part, and this proves the theorem.**

Note: **Same dovetailing explanation holds good for d/dt (partial differential with respect to t), d/dt,d^2/dt^2 (acceralation: double dot). Such an exposition is helpful in optimisation problems, duality of motion, wave motion a homogeneous functions studies, and control theory which we intend to incorporate in future.** It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## Persistence, Coexistence And Thrusty Permanence For Interacting Structured Particles, Transactions In Finance, Econometric Productivity Functions, Environment, Interfaces, Cosmological Interactions Et Al

## Josef Hofbauer And Sebastian J. Schreiber Model Thereof: A Functional Appraisal

ABSTRACT. **JOSEF HOFBAUER AND SEBASTIAN J. SCHREIBER** studied the dynamics of interacting structured systems (in their case population dynamics) can be modeled by $\frac{dxi}{dt} = A_i\ (x)\ xi$ where $x_i\ \epsilon\ R^n_{i,}\ x = (x_1, \ldots, x_k)$, and $A_i(x)$ are matrices with non-negative off- diagonal entries. These models are permanent if there exists a positive global attractor and are robustly permanent if they remain permanent following perturbations of $A_i\ (x)$. Necessary and sufficient conditions for robust permanence are derived using dominant Lyapunov exponents $X_i\ (\mu)$ of the $A_i\ (x)$ with respect to invariant measures $\mu$. The necessary conditions requires $\max_i\ X_i(\mu)>0$ for all ergodic measures with support in the boundary of the non-negative cone. The sufficient condition requires that the boundary admits a Morse decomposition such that $\max_i$ Xi$(\mu)>0$ for all invariant measures $\mu$ supported by a component of the Morse decomposition. When the Morse components are Axiom A, uniquely ergodic, or support all but one population, the necessary and sufficient conditions are equivalent. Applications to spatial ecology, epidemiology, and gene networks are given. Appeared in *Journal of Differential Equations*, 248, 1955-1971 (2010) It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model

becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## INTRODUCTION

A fundamental issue in population biology is what are the minimal conditions to ensure the long-term survivorship for interacting populations whether they be viral particles, bio-chemicals, plants, or animals. When these conditions are met the interacting populations are said to persist or coexist. Since the pioneering work of Lotka and Volterra on competitive and predator prey interactions, Thompson, Nicholson, and Bailey on host-parasite interactions, and Kermack and McKendrick on disease outbreaks, nonlinear difference and differential equations have been used to understand conditions for population persistence [1-5]. One particularly important form of persistence for deterministic models is *permanence* or *uniform persistence* which corresponds to the existence of a global attractor bounded away from extinction of one or more species. When such an attractor exists, interacting populations are able to recover from "vigorous shake ups" of the of the population state. Permanence has been characterized from a topological perspective [6-8] and with average Lyapunov functions [9-12] for differential equation models of the form

$$\frac{dxi}{dt} = x_i f_i(x) \qquad i = 1, \ldots, k \qquad (2)$$

Where $x_i$ is the abundance of population $i$, $x = (x_1, \ldots, x_k)$, and $f_i(x)$ is the per-capita growth of population $i$.

An important extension of the concept of the permanence is *robust permanence* [13-16] which requires that (1) remain permanent after sufficiently small perturbations of the per-capita growth rates $f_i$. The importance of robust permanence as one crude form of structural stability for population models are approximations to reality. Consequently, if nearby models (e.g. more realistic models) are not permanent despite the focal model being permanent, and then one can draw few conclusions about the persistence of the biological system being approximated by the model. One can view robust permanence as one crude form of structural stability for debits-

credits particles, propositions, transactions in the celestial subterranean realm and ceratoid dualism, finance like debiting say a current account and crediting to another account, or recrediting to the account which had been earlier debited, further conduction of the cash credit account, black hole formation, guzzling of inventories for the output, creation of income from output, reinvestment of income to investment,gormadizing of matter by black holes, and emission of radiation therefore, gamma ray bursts, in fact most of the transactions and interactions in the universe are conservative and quantum mechanical universe, something like a account acting as predator and prey. More often than not as we have seen from Haimovici model such models could be applied to the problems, propositions, which are not subsumed by the concomitant problems in the bearing. the population models. For this perspective, it is not unexpected that there are permanent systems that cannot be approximated by robustly permanent systems [17]. In [14, 16], criteria for robust permanence were developed with respect to the average per-capita growth rates $\int f_i(x)d\mu(x)$ with respect to invariant probability measures $\mu$ supported on the boundary of the positive orthant (i.e. where one or more Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates are extinct). Roughly, these criteria for robust permanence require that the average per-capita growth rate (i.e. invasion rate) is positive for some missing species.

While (1) can account for many types of Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates. Interactions, it assumes that all individual within populations are exactly the same. However, theoretical biologists have long recognized that different individuals within a population may be in different states (e.g. different sizes or ages, living in different spatial locations) and these differences can have important consequences for population dynamics [18-24]. To account for how these differences influence persistence of interacting populations, we develop criteria for robust permanence for the dynamics of $k$ interacting structured populations. Roughly, these dynamics correspond to replacing $x_i$ in (1) by vectors and $f_i(x)$ by matrices. In section 2, we describe these models in greater detail and review some basic concepts from dynamical systems theory. In section 3, we introduce the structured analog of average per-capita growth rates and define the concept of robustly unsaturated invariant sets. Necessary and sufficient conditions for an invariant set to be robustly unsaturated are proven. For Axiom A or uniquely ergodic invariant sets, the necessary and sufficient criteria are shown to be equivalent. In section 4, we use Morse decompositions and the criteria for robustly unsaturated invariant sets to develop necessary and sufficient conditions

for robust permanence. In sections 5 and 6, we provide several applications of our results to spatially structured ecological models, structured epidemiological models, and models of gene networks. We apply the model to econometrics. To give an example investment needs inventory to be used towards the consummation of output, part of output is to be sold for consolidation of income, and some income might be replenished in the investment and so on. In finance a cash credit account could be debited or credited when there is say purchase of demand bills or discount of usance bills. Cash credit account is debited whenever cheques are issued towards the purchase of raw materials and so on. In fact some of the models have been applied to Organic psychiatry and we have shown how such models are useful in dealing with the propositions of Husserl, Derrida, Deleuze and others.

## 2. MODELS AND ASSUMPTIONS

Let $x_i$ denote the state of the $i$-th population that lies in the non-negative cone $C_i$ of $R^{ni.}$ Define $x=(x_1,\ldots x_k)\epsilon C$ to be the non-negative cone of $R^n$ where $n=\sum_{i=1}^{k} n_i$. Let $C_i^+$ denote positive cone { $x_i \epsilon\ C_i$: $\prod_i$i noitalupop rof$^j$

$i$ *and* $C^+=\prod_i C^+_i$ , the positive cone for the interacting populations. If $x \mapsto A_i\ (x)$ is a map into $n_i x n_i$ matrices that describes the growth of population I, then the dynamics of the interacting populations are given by

$$(2)\frac{dx_i}{dt} = A_i(x)x_i \ i=1,\ldots,k$$

Let $x.t$ denote the solution of (2), with the initial condition x, and more generally the semi flow generated by (2). Before stating our assumption on (2), recall a few definitions from dynamical systems. Given sets $I \subseteq R$ and $K \subseteq C$, let $K.I = \{x.t : t \epsilon I, x \epsilon K\}$. A set $K \subseteq C$ *is invariant* if K.t = K for all t>0, and *forward invariant* if $K.t \subset K$ for all $t > 0$. The *omega limit* set of a set $K \subseteq C$ equals $\omega(K) = \bigcap_{t\geq 0,} \overline{K.[t,\infty)}$. The *alpha limit* set of $K \subseteq C$ is $a(K) = \bigcap_{t\leq 0} \overline{K.(-\infty,t]}$. Given a forward invariant set K, $B \subset K$ is an attractor for the semi flow $x.t$ *restricted* to K provided there exists an open neighborhood $U \subseteq K$ of B such that $\omega\ (U) = B$. The stable set of a *compact invariant* set $K$ is defined by

$$W^s(K) = \{x \epsilon C :\ \omega\ (x) \neq \emptyset \text{ and } \omega\ (x) \subseteq K\}.$$

The semi flow generated by (2) is *dissipative* if there exists an attractor B with $W^s\ (B) = C$.

Throughout this paper, we make the following assumptions:

**A1**: $x \mapsto A_i(x)$ is continuous

**A2:** $A_i(x)$ is irreducible and has non-negative off- diagonal entries,

**A3:** $x.t$ is defined for all $t \geq 0$

**A4:** $x.t$ is dissipative with global attractor $\Gamma(A)$.

A1 is a basic regularity assumption. The non-negative of off- diagonal entries in A2 implies that there are no negative feedbacks between individuals of different states in population *i*. This assumption is meet for many types of structured models as discussed in sections 5 and 6. The irreducibility assumption A2 is generally meet and implies that all individuals within a population can pass through all states. In Remark1, we discuss how this irreducibility assumption can be relaxed. A3 ensures that the population dynamics are defined for all future time. A4 requires that population densities/ abundances eventually are uniformly bounded, a condition that should be meet for an biologically realistic model.

We say that the semi flow of (2) is *permanent* if this semi flow is dissipative and there is a positive attractor $B \subset C^+$ such that $W^s(B) \subset C^+$. We will denote this positive attractor by $\Lambda(A)$. Permanence ensures that populations can recover from rare large perturbations and allows for a diversity of dynamical behaviors.

For K ⊂ C and $\delta > 0$, define the $\delta$- *neighborhood* of K as

$$N\delta(K) = \{x \in C : \backslash x - y < \delta \text{ for some } y \in K\}.$$

We define a $\delta$- *perturbation of (2)* to be a system of the form

$$\frac{dx}{dt} = \tilde{A}(x)x$$

that satisfies assumptions A1-A4, $\Gamma(\tilde{A}) \subset N(\Gamma(A))$, *and* $||A(x) - \tilde{A}(x)|| \leq \delta$ *for all* $x \in N(\Gamma(A))$. (2) is *robustly permanent* if there exist >0 and $\epsilon > 0$ such that all $\delta$- perturbations of (2) are permanent, and $d(\Lambda(\tilde{A}), C\backslash C^+) > \epsilon$ for all $\delta$- perturbations i.e., there is a common/ uniform region of repulsion around the boundary.

## INUNDATION OF COMPACT SETS

We begin by studying the linear skew product flows on $\Gamma(A) \times R^n_i$ defined by $(x,y).t = (x.t, B_i(t,x)y)$ where $Y(t) = B_i(t,x)$ is the solution to $\frac{dY}{dt}(t) = A_i(x.t)Y(t)$ with y(0) equal to the identify matrix. Our assumption that $A_i$ is irreducible with non- negative off diagonal entries implies that $B_i$ (t,x)$C_i \subset C_i^+$ for all $x$ and t > 0 (see, e.g., [25] ). A

result of Ruelle [26, Prop, 3.2] implies that there exist continuous maps $u_i\, v_i : \Gamma(A) \rightarrow C_i^+$ with $u_i(x)|=|v_i(x)|=1$ such that The line bundle $E_i(x)$ spanned by $u_i(x)$ is invariant i.e. $E_i(x.t)= B_i(x,t)E_i(x)$ for all t ≥ 0.The vector bundle $F_i(x)$ perpendicular to $v_i(x)$ is invariant i.e. $F_i(x.t)= B_i(x,t)F_i(x)$ for all $t \geq 0$.There exist constants $a> 0$ and $\beta>0$ such that(3) $||B_i(t,x)F_i(x)|| \leq a\, exp(-\beta t)||B_i(t,x)E_i(x)||$ For all $x \in K$ and $t \geq 0$.**Interaction rates.** Given any x $\epsilon$ $\Gamma$(A), define the *interaction and inundation rate* of *species i at particle/proposition/problem/ state x as* $\lambda_i(x)= limsup_{t\rightarrow\infty}\left(\frac{1}{t}\right)$ *in* $||B_i(t,x)||$ when $x_i$ =0, $\lambda_i(x)$ provides an upper bound to the rate of growth of population *i* when introduced at infinitesimally small densities. Important properties of this invasion rate are summarized in the following two propositions. For instance, the second proposition implies if $(x.t)_i$ stays bounded away from zero (i.e. population *i* persist), then

$\lambda_i(x)=0$. Alternatively, if $\lambda_i(x)<0$,then population *i* is doomed to extinction. **Proposition 1.** $\lambda_i(x)$ *satisfies the following properties:* $\lambda_i(x)= limsup_{t\rightarrow\infty}\left(\frac{1}{t}\right) In \quad |B_i \quad (t,x)v|$ *for any* $v>0$ *in* $R^{ni}$, *and* $\lambda_i(x)= limsup_{t\rightarrow\infty}\left(\frac{1}{t}\right) \int^t_0 (A_i(x.s)u_i(x.s),u_i(x.s))ds.$

**Proof:** To prove the first property, we first show that

$$\lambda i(x) = limsup_{t\rightarrow\infty} \frac{1}{t} In\ |B_i(t,x)u_i(x)|.$$

To this end, let $v \in R^{ni}$ be any non-zero vector. Since $R^{ni}= E_{i(x)} \quad F_i(x)$, we can write $v=au_i(x)+w$ with $a \in R$ and $w \in F_i(x)$. Equation (3) implies

$$|B_i(t,x)v| \leq a|\, B_i(t,x)u_i(x)|+ |B_i(t,x)w| \leq |\, B_i(t,x)u_i(x)|(a+ \alpha+ exp\ (-\beta t)|\, w|).$$ Hence, $limsup_{t\rightarrow\infty} \frac{1}{t} In\ B_i(t,x)\ v| \leq$

$$limsup_{t\rightarrow\infty} \frac{1}{t} In\ B_i\,(t,x)\ u_i\,(x) \leq \lambda_i(x)$$

for all non zero vectors $v \in R^{ni}$. Since $||B_i(t,x)|| = sup|v|=1B_i(t,x)v|$, this inequality implies that $\lambda i(x) = \limsup_{t\rightarrow\infty} \frac{1}{t}$ In $|B_i(t,x)v|$, this inequality implies that $\lambda_i(x) = limsup_{t\rightarrow\infty}\left(\frac{1}{t}\right)$ In $|t,x)u_i(x)|$. Now let v $\epsilon$ $C_i^+$. Then, we can write $v = au_i(x)+ w$ with $a$ >0 and $w \in F_i$.

Equation (3) implies

$$|B_i(t,x)v| \geq a|B_i(t,x)u_i(x)|-|B_i(t,x)w|$$

$$\geq |\, B_i(t,x)u_i(x)|(a- \alpha+ exp\ (-\beta t)|\, w|).$$

Since a>0, this inequality implies that

$$\lambda i(x) \geq \limsup_{t\to\infty} \frac{1}{t} \text{In } |B_i(t,x)v| \geq \lambda i(x) \geq \limsup_{t\to\infty} \frac{1}{t} \text{In } |B_i(t,x)u_i(x)| = \lambda i(x)$$

To prove the second assertion, let $b_i(t,x) = In\ |B_i(t,x)u_i(x)|$. Invariance of $u_i(x)$ implies that $b_i(t,x)$ is additive:

$$B_i(t+s,x) = In|B_i(t+s,x)\ u_i(x)| = In\ |B_i(s,x.t)B_i(t,x)u_i(x)$$

$$= \text{In } |B_i(s,x.t)\ u_i(x.t)||\ B_i(t,x)\ u_i(x)|$$

$$= \text{In } b_i(s,t.x)\ b_i(t,x)$$

Additivity of $b_i(t,x)$ and the fact that $b_i(0,x)$ =0 implies

$$\frac{d}{dt} b_i(t,x) = \lim_{s\to 0} \frac{bi(t+s,x-bi(t,x)}{S-}$$

$$= \lim_{s\to 0} \frac{bi(s,x.t)}{S} = \frac{d}{ds}|_{s=0}\ b_i(s,x.t)$$

=

$$=(A_i(x.t)u_i(x.s),u_i(x.s),$$

The Fundamental Theorem of Calculus implies

$$B_i(t,x)= \int_0^t (A_i(x.s)\ u_i(x.s),\ u_i(x.s))ds\backslash$$

and the second assertion follows.

**Proposition 2:** *For the solutions of (2) we have:* if $x_i>0$, then $\lambda i(x) \leq 0$. If $\lambda i(x)<0$, then $\lim_{t\to\infty} (x.t)_i=0$ If $limsup_{t\to\infty}$ $(x.t)_i|>0$, then $\lambda i(x)=0$.

**Proof:** First, assume that x satisfies $x_i>0$. Since the semi flow is dissipative, there exists $\gamma >0$ such that $|x.t| \leq \gamma$ for all $t \geq 0$. Proposition 1 and the definition of the skew product flow imply

$$\lambda i(x) = \limsup_{t\to\infty} \frac{1}{t} \text{In } |B_i(t,x)x_i| = \limsup_{t\to\infty} \frac{1}{t} \text{In } |\ (x.t)_i\ |$$

$$\leq \limsup_{t\to\infty} \frac{In\ \gamma}{t} = 0$$

To prove the second assertion, assume that $\lambda i(x) < 0$. *If* $x_i$=0, then the invariance of the faces of C imply that $(x.t)_i$=0 for all t≥ 0. Alternatively, if $x_i$> 0, then Proposition I and the skew product flow imply

$$\limsup_{t\to\infty} \frac{1}{t} In|(x.t)_i| = \lambda i(x) < 0.$$

The final assertion follows from the first two assertions. 5

**3.2. Invariant measures.** We review some definitions from ergodic theory. Given a Borel probability measure $\mu$ on ***C***, the *support* of $\mu$, denoted supp( $\mu$),is the smallest closed set whose complement has measure zero. A Borel probability measure $\mu$ is called *invariant* for (2) provided that Integral $h(x)d\mu(x)$ = Integral $h(x.t)d\,\mu(x)$for all $t \geq 0$ and for all bounded continuous functions $h$:C→R. An invariant measure $\mu$ is called *ergodic* provided that $\mu(B)$=0 or 1 for any invariant Borel set *B.*

For an invariant measure $\mu$ define the invasion rate species $i$ with respect to $\mu$ as

$$\lambda_i(\mu) = \int_c \langle A_i(x)u_i(x), u_i(x)\rangle d\mu(x)$$

**Proposition 3:** Let $\mu$ be an *invariant measure.* $\lambda_i(\mu)$ *Satisfies the following properties:*$\lambda_i(x) = \lim_{t\to\infty} \frac{1}{t}\int_0^t \langle A_i(x.s)u_i(x.s), u_i(x.s)\rangle, ds$ exists for $\mu$ *- almost every* $x$. *Moreover, if* $\mu$ *is ergodic,* $\lambda_i(x) = \lambda_i(\mu)$ for $\mu$-*almost every* $x$.*If* $\mu$ *is ergodic, then there exists* $I \subset \{1, \dots, k\}$such that $\mu(\Pi_{i\in I} \boldsymbol{C}_i^+)$=1 and $\lambda_i(\mu) = 0\ for\ all\ \ i \in I$.

**Proof:** Let $\mu$ be an invariant measure. Define $h_i(x) = \langle A_i(x)u_i(x), u_i(x)\rangle$. The Birkhoff's' ergodic theorem and Proposition 1 imply

$$\lambda_i(x) = \lim_{t\to\infty} \frac{1}{t}\int_0^t h_i(x.s)ds$$

Exists for $\mu$- almost every $x$. Moreover, $\lambda_i(x) = \lambda_i(\mu)$ $\mu$- almost surely if $\mu$ is ergodic.

Assume $\mu$ is ergodic. By ergodicity and invariance of the faces of **C**, there exists $I \subset \{1, \dots, k\}$ such that $\mu(\Pi_{i\in I} \boldsymbol{C}_i^+)$=1. Let K $\subset \Pi_{i\in I} \boldsymbol{C}_i^+$ be a compact set such that $\mu(K) > 0$. The Poincare recurrence theorem and

the Birkhoff ergodic theorem imply that there is $x \in \Pi_{i\in I} \boldsymbol{C}_i^+$ such that $x.t_n \in K$ for some $t_n \uparrow \infty$ and $\lambda_i(x) = \lambda_i(\mu)$ for all $i$. The third assertion of Proposition 2 implies $\lambda_i(\mu) = 0$ for all $i \in I$.

**3.3. Robustly unsaturated sets**. Let $K \subset \partial C$ be a compact isolated invariant set for the flow $x.t$ restricted to $\partial C$. $K$ is unsaturated if $W^s(K) \subset \partial C$ and $K$ is isolated for $x.t$. If $K$ is not unsaturated, then $K$ is saturated. $K$ is robustly unsaturated for (2) if there exists $\delta > 0$ such that the continuation of $K$ for any $\delta -$ perturbation of (2) is unsaturated.**Theorem 1:** Let $K$ be a compact isolated invariant set for $x.t$ restricted to$\partial C$. If one of the following equivalent conditions hold For all invariant measures $\mu$ supported by $K$ (4) $\max_{1\le i\le k} \lambda_i(\mu) = 0$

There exist $p_i > 0$ such that (5) $\sum_{1\le i\le k} p_i \lambda_i(\mu) = 0$

For all ergodic probability measures supported by $K$.

then $K$ is robustly unsaturated for (2). Alternatively, if $x \to A(x)$ is twice continuously differentiable and $K$ is robustly unsaturated, then (4) holds for all ergodic measures $\mu$ supported by $K$.

**Remark 1**: For some applications (e.g., the disease model considered in section 6.1) it useful to relax the irreducibility assumption A2. For instance, if there exists an open neighborhood U of K such that for each $i$, $A_i(x)$ has a fixed off-diagonal sign pattern for all $x \in U$,then $A_i(x)$ can be decomposed into a finite number, say $m_i$, of irreducible components. For each of these irreducible components, one can define $\lambda_i^j(\mu) = \int_C \langle A_i^j(x) u_i^j(x), u_i^j(x)\rangle \, d\mu(x)$where $A_i^j(x)$ is the submatrix of $A_i(x)$ corresponding to the $j$-th irreducible component of $A_i(x)$ $and$ $u_i^j(x)$ is the continuous invariant sub bundle for the irreducible linear cocycle determined by $A_i^j(x)$. If we define $\lambda_i(\mu) = \max_{1\le j\le m_i} \lambda_i^j(\mu)$, then all of the assertions of Theorem 1 still hold.

**Proof:** To prove the first assertion of the theorem, we need the following lemma. We call an invariant measure $\mu$ for (2) saturated if $\lambda_i(\mu) \le 0$ for all $i$

**Lemma 1**: Let $\delta_n$ be a non –negative sequence that converges to zero as n$\to\infty$. If $\mu_n$ are saturated invariant measures for $\delta_n -$ perturbations of (2), then the weak$^*$ limit points of $\{\mu_n\}_{n=1}^{\infty}$ is a non-empty set consisting of saturated invariant measures for (2).

**Proof:** Let $\delta_n$ be a non-negative sequence that coverages to zero. Let $\mu_n$ be saturated invariant measures for $\delta_n$ perturbations, $\dot{x} = A^n(x)x$, of $\dot{x} = A(x)x$. By weak$^*$ compactness of Borel probability measures supported on $N_1(\Gamma(A))$, there exist weak$^*$ limit points of $\mu_n$ . Let $\mu$ be such a weak$^*$ limit point. Since $\mu_n$ are supported by $N_1(\Gamma(A))$, for all $n$ sufficiently large, $\mu$ is supported by $N_1(\Gamma(A))$.To verify that $\mu$ is an invariant measure for (2),let $h: C \to R$ be a bounded continuous function. Let $t > 0$ and $\in > 0$ be given. Choose $n$ sufficiently large so that $|\int_c h(x)\, d\mu(x) - \int_c h(x) d\mu_n(x)| \leq \in$, $|\int_c h(x.t)\, d\mu(x) - \int_c h(x.t) d\mu_n(x)| \leq \in$, and $|h(x.t) - h(x^n.t)| \leq \in$ for all $x \in N_1(\Gamma(A))$.Then $|\int_c h(x.t) - h(x) d\mu(x)|$ is

$\leq \quad |\int_c h(x.t)\, d\mu(x) - \int_c h(x.t)\, d\mu_n(x)| + |\int_c h(x.t)\, d\mu_n(x) - \int_c h(x)\, d\mu(x)|$

$\leq \quad \in + \int_c |h(x.t) - h(x^n.t)|\, d\mu_n(x)| + |\int_c h(x^n.t)|\, d\mu_n(x) - \int_c h(x) d\mu(x)|$

$\leq \quad 2 \in + |\int_c h(x^n.t) - h(x)\, d\mu_n(x)| + |\int_c h(x) d\mu_n(x) - \int_c h(x) d\mu(x)|$

$\leq \quad 3 \in$

Where the last line follows from the invariance of $\mu_n$ for $x^n.t$. Since $\in > 0$ and t > 0 are arbitrary, we have $\int_c h(x.t) d\mu(x) = \int_c h(x) d\mu(x)$ for all t >0 and all bounded continuous functions $h: C \to R$. It follows that $\mu$ is an invariant measure for (2). To see that $\mu$ is saturated, define $h_i^n(x) = \langle A_i^n(x) u_i^n(x), u_i^n(x) \rangle$ where $u_i^n(x)$ spans the invariant one-dimensional bundle given by $[26, Prop. 3.2]$ for $\frac{dx}{dt} = A^n(x)x$. Since $A_i^n(x) \to A_i(x)$ uniformly for $x \in N_1((A))$ as $n \to \infty$, $[26, Prop. 3.2]$ implies that $u_i^n(x) \to u_i(x)$ converges uniformly for $x \in N_1(\Lambda(A))$. Given $\in > 0$, choose $n$ sufficiently large so that $|h_i^n(x) - h_i(x)| \leq \in$ for all $x \in N_1(\Lambda(A))$ and $|\int_c h_i^n(x)\, d\mu(x) - \int_c h_i^n(x)\, d\mu_n(x)| \leq \in$.

Then

$$\lambda_i(\mu) \leq \int_c |h_i(x) - h_i^n(x)| d\mu(x) + |\int_c h_i^n(x)\, d\mu(x) - \int_c h_i^n(x)\, d\mu_n(x)|$$

$+ \int_c h_i^n(x)\, d\mu_n(x)$

$\leq 2\epsilon$

Since $\epsilon > 0$ is arbitrary, $\lambda_i(\mu) \leq 0$ for all $i$ and $\mu$ is saturated.

Using this lemma, we prove that if K is saturated, then there exists a saturated invariant measure $\mu$ supported by K. Assume K is saturated. Work of Hofbaue and So $[7, Thm.\ 2.1]$ implies that either there exists $\mathcal{Y} \in W^s(K) \cap C^+$ or K is not isolated for the unrestricted flow $x.t$ If there exists $\mathcal{Y} \in W^s(K) \cap C^+$, then for all $t \leq 0$ define $v_t = \frac{1}{t} \int_0^t \delta_{y.s}\, ds$ where $\delta_{y.s}$ denotes a Dirac measure based at the point $\mathcal{Y}.s$. Dissipativeness of (2) and weak* compactness of the Borel probability measures supported by (A) imply there exist $t_k \to \infty$ such that the sequence $v_{tk}$ converges in the weak* topology to a Borel probability measure $\mu$ with support in K.A standard argument implies that $\mu$ is $x.t$ invariant. Define $h_i: C \to R$ by $h_i(x) = (A_i(x)u_i(x), u_i(x))$.Proposition 2 and weak* convergence imply that

$$0 \geq (\lambda i(y) = \limsup_{t\to\infty} \frac{I}{t} \int_0^t h_i(y.s)ds$$

$$= \limsup_{t\to\infty} \int_c h_i\, dv_t$$

$$\geq \limsup_{t\to\infty} \int_p h_i\, d\, v_{tk} = \lambda i(\mu)$$

For all $i$. Alternatively suppose that K is not isolated for the semi flow Then there exists a sequence of positive $\omega - limit$ sets that accumulate on K. Let $\mu_n$ be a sequence of ergodic probability measures supported by these $\omega -$ limit sets. Propositions 1 and 3 imply that $\mu_n$ are saturated for all $n$ Applying Lemma I with $\delta_n$=0 for all $n$ implies that there exists a saturated invariant measure $\mu$ suppoorted by *K*. Hence we have shown that if *K* is saturated then there exists a saturated invariant probability measure $\mu$ supported by *K* Equivalently, (4) holding for all invariant probability measures with support in *K* implies that *K* is unsaturated.

Next we show that if *K* not robustly saturated, then there exists a saturated invariant measure supported by *K* Indeed suppose *K* is not robustly saturated, Then there exists a non-negative sequence $\delta_n$ converging to zero and a

sequence of saturated measures $\mu_n$ for $\delta_n$-perturbations of (2)with support in the continuation of *K*. Let $\mu$ be a weak* limit point of $\{\mu_n\}\infty_{n=1}$ Lemma 1 implies that $\mu$ is saturated. Moreover, since the continuation of *K* converges to *K* as $\delta_n \to 0, \mu$ is a saturated invariant measure for (2) supported by *K* Hence we have shown that if *K* is not robustly saturated then there exists a saturated invariant measure supported by *K* Equivalently if (4) holds for all invariant measures with support in *K* then *K* is robustly unsaturated.

To see the equivalence of the conditions given by (4) and (5) let $\Delta = \{p \in R_+^k : \sum_i P_i = 1\}$ and notice that

$$\text{Min max } \lambda i\ (\mu = \min_{\mu}\ \max_{p\epsilon\Delta}\ \sum_i\ Pi\lambda i(\mu)$$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K* The Minimax theorem (see, e. g.,[27] implies that

$$(6)\ \min \max\ \lambda i(\mu) = \min_{\mu}\ \max_{p\epsilon\Delta}\ \sum_i\ Pi\lambda i(\mu)$$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K.* Since min$\mu$ $\sum_i\ P_i\ \lambda_i(\mu)$ is attained at an ergodic probability measure with support in *K*, the equivalence of the conditions given by (4) and (5) is established.

To prove the final assertion of the Theorem assume $x \to A(x)$ is twice continuously differentiable and there exists a saturated ergodic measure $\mu$ supported by *K* Proposition 3 implies that there exists *I*C {1,…,k} such that (□$i \in iC_i^+$)=1and $\lambda i(\mu) = 0$ For all $i \in I. SinceKc\ .\ \partial$□,$\{1, \ldots . k\}\backslash I$ is non-empty. We will show that (2) is not robustly permanent by proving that for all $\delta > 0$ there exists a $\delta$ −perturbation of (2) for which the continuation of *K* is saturated Let $\delta > 0$ be given Choose V $c$ N$\delta$ ($\Gamma$(A)) to be a compact neighborhood of Γ(A) such that V.t $c$ int *V* for all t > 0.Let $\eta > 0$ be such that *W*=:*N* $\eta$(Γ(A))$c$ *V* Let p:c→ [0,1]$be$ a$C^\infty$function such that $p(x) = 1 \quad for\ all\ x \in \Gamma(A) and\ p(x) = 0 for\ all\ x \in C\backslash W\ Define\ \tilde{A} = (\tilde{A}_1, \ldots, \tilde{A}_k)$ by

$$\tilde{A}_i(x) = \{\frac{A(x)}{A(x)-\frac{\delta}{2Idp\,(x)}} \quad \frac{if\ i\epsilon I}{if\ i \notin I}$$

Where Id denotes the identity matrix of appropriate dimension Let x. t denote the semi flow of $x = \tilde{A}(x)x$.Since$x.t = x.t$ $whenever\ x.[0,t] \in C\backslash W, it\ follows\ that\ \ \tilde{V}.t \subset intV\ for\ all\ t >$ $0. Hence, \Gamma(\tilde{A}) c\ V\ c\ N\delta\ (\ (A))). We\ \ also\ have\ ||A(x) - \tilde{A}(x)|| \leq \frac{\delta}{2}$Therefore $x = \tilde{A}(x) is\ a\ \delta -$ perturbation of (2) By construcation,$x.t = x.t\ for\ all\ x \in K and\ t \geq 0. Hence, \mu$ is ergodic for the semiflow of $\dot{x} = \tilde{A}(x)x\ and\ \lambda_i(\mu) \leq -\frac{\delta}{2}$ for this semiflow and $i \in \{1, \dots, k\}\backslash I$.Let $\mathcal{L}$ and O be the Lyapunov exponents and Oseledec regular points supported by $\mu$ for $x = \tilde{A}(x)x(see, e.g.,$ [28] for definitions).At each point $x \in O$, the splitting of $R^n$ determines three subspaces: the stable subspace $E^s(x)$, the center subspace $E^c(x)$, and the unstable subspace $E^u(x)$ Proposition 3 and our choice of Ã imply that $E^s(x) \cap C^+ \neq \emptyset$. The Pesin stable manifold theorem [28, Corollaries 3.17 and 3.18] implies that tangent to $E^s(x)$ $E^c(x)$ and $E^u(x)$ are locally $x.t-$invariant families of $C^1$ discs $W^s_x, W^c_x$ and $W^u_x$ corresponding to the stable ,center and unstable manifolds The family of stable manifolds $W^s_x$ is contained in $W^s$ (Ã,supp $(\mu)$). $Since\ E^s(x) \cap C^+ \neq \emptyset$, $W^s_x \cap C^+ \neq \emptyset$ for some $x \in K$. Consequently $K$ is saturated for $x = \tilde{A}(x)x$.

The difference between the sufficient and necessary condition is that for the sufficient condition (4) has to hold for all invariant measures supported by *K*, while for the necessary condition (4) has to hold only for ergodic measures supported by *K*. since the invariant measures lie in the convex hull of the ergodic measures the sufficient condition can be more restrictive then the necessary condition.

For three Classes of invariant sets the necessary and sufficient conditions coincide Recall a compact invariant set *K* for (2) is uniquely ergodic if K only supports one invariant measure Recall a compact invariant set *K* is Axiom A if the flow of (2) restricted to *K* is transitive and hyperbolic (see, e.g.,[29] for definitions.)

**Corollary 1:** Assume $x \to A(x)$ is twice continuously differentiable It K$\subset \partial$□ is a compact invariant set and either is uniquely ergodic or Axiom A, then *K* is robustly unsaturated if and only if (4) holds for all ergodic, measures $\mu$ with support in *K*

Proof. If K is uniquely ergodic then the assertion follows immediately For an Axiom A invariant set Sigmund [30, thm.1] has proven that invariant measures supported by periodic orbits of $K$ are dense in the set of invariant probability measures supported by $K$ In particular any invariant measure can be approximated by an ergodic measure and consequently the result follows.

Another special case where the necessary and sufficient conditions coincide is when $K$ supports all populations except one for discrete time models this case was considered by [22].

**Corollary2:** Assume $x \to A(x)$is twice continuously differentiable Let K$c\ II_i \geq 2C_i^+$ be a compact invariant set. Then K is robustly unsaturated if and only if $\lambda_1(\mu) > 0$ for all invariant measures supported by $K$

## 4. MORES DECOMPOSITIONS AND ROBUST PERMANENCE

To state the sufficient condition for robust permanence, we use a characterization of permanence due to Garay [8] and Hofbauer and So [7] that involves Morse decompositions of the boundary flow Conley[31] defined a collection of sets$\{M_{1,\ldots,}M_m\}$ to be a Morse decomposition for a compact invariant set $K$ if $M_1\ldots,M_m$ are pairwise disjoint, compact isolated invariant sent for the flow of (2) restricted to K. For each $x \in K$ there are integers r=r$(x) \leq s = s(x)$ such that $a(x) \subseteq M_r$and $\omega(x) \subseteq M_{s.}$ If $r(x) = s(x), then\ x \in M_{r.}$ Garay, Hofbauer and so [7,8] proved the following characterization of permamnce

**Theorem 2** (Garay, Hofbauer-so) if $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$then (2) is permanent if and only if each of the components $M_1$ are unsaturated.

Theorem 1 and 2 imply the following result: As stated in the above all these results are true for most of problems, interactions, transactions in cosmology, econometrics, finance and other subjects, which we shall illustrate as we go on constructing models for the same.

**Theorem 3:** If $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$and (4) holds for each of the components of the More decomposition, then (2) is robustly permanent Conversely if $x \to A(x) is\ twice\ continuously\ differentiable$

*and (4) is violated by an ergodic measure supported by one of the components of the Morse decomposition, then (2) is not robustly permanent.* Theorem 3 in conjunction with Corollaries 1 and 2 yield a characterization of robust permanence for a class of structured models.

**Corollary 3:** *Let $\{M_1,.....,M_m\}$ os a Morse decomposition for $\partial C$. Assume $x \mapsto A(x)$ is twise continuously differentiable and for each Morse component $M_i$ one of the following assertions hold $M_i$ is Axiom A, $M_i$ is uniquely ergodic, or There exists $j \in \{1,....,k\}$ such that $M_i$ ----- $\{ x \in C : x_l \gg 0$ for $l \neq j\}$ Then (2) is robustly permanent if and only if (4) holds for all ergodic measures supported by $U_i M_i$.*

Our result also provided a structured analogue for characterizing totally permanent systems[32]. For an ergodic probability measure$\mu$, define sup ($\mu$) to be the subset $I$ -----$\{1,\ldots,k\}$ such that $\mu(---_{i \in I} C_+) = 1$

**Corollary 4:** *The following statements are equivalent: For all ergodic probability measures $\mu$ with support in $\partial C$, $\lambda_i(\mu) > 0$ for all $I \in \{1,...,k\} \setminus supp(\mu)$ (2) and all of its subsystems are robustly permanent.*

*Proof. Suppose the first statement holds. Then unsaturated condition given by (5) with p= (1,...,1) and Theorem 3 implies that (2) and all its subsystems are robustly permanent.. The other direction follows form the necessary condition for robust permanence in Theorem 3 and the second assertion of Proposition 3.*

*To illustrate the broad applicability of theorem 3, we develop applications to spatially structured models in section 5 and to a disease model and a gene network model in section 6.*

## 5. APPLICATIONS: PATCH MODELS

A fundamental application of our result are to spatially structured models, with *k* species dispersing between *m* patches {33-35}:

(7) $$\frac{dx_i^j}{dt} = x_i^j f_i^j(x^j) + \sum d_i^{jl} x_i^l - e_i^j x_i^j \qquad i = 1,..,k \qquad j=1,\ldots,m$$

*l*

Where $x_i^j$ denotes the density of species I in patch $j$, $x^j = (x^j_1, \ldots, x^j_k)$ is the vector of species densities in patch $j$, $f^j_i$ is the per-capita growth rate of species $i$ in patch $j$, $d_i^{jl} \geq 0$ is the dispersal rate for species $i$ from patch $l$ into patch $j$, and $e_i^j$ is the emigration rate of species $i$ out of patch $j$. Hence $e_i^j \geq \sum_l d_i^{lj}$. We assume that the matrices $(d_i^{lj})_{j,l}$ are irreducible for each species $i$. it is then easy to write (7) in the form (2), with $x_i = (x^1_i, \ldots, x_i^m)$ 'where' denotes transpose.

For a single species (7) generates a monotone flow. Under mild assumption, e.g., each $f^j_i$ *is* decreasing and negative for large densities, there is a globally stable equilibrium, see, e.g.,[36] and [37, sect. 5.4]. This equilibrium will be positive if the invasion rate at the origin $\lambda(0)$, which is given by the leading eigenvalue (stability modulus) of the matrix $(d^{jl} + f^j(0) - e^j)$, is positive.

5.1 Two types of debitable and creditable systems:. For *two competing species,* e.g., $f^j_i$ is decreasing with respect to both species and is negative at large densities, the dynamics are still monotone, and hence almost all orbits converge to an equilibrium. Robust permanence requires two conditions. First, $\lambda_i(0) > 0$ for both species in which case there are two single species equilibria $E_1$ and $E_2$. Second, the invasion rates $\lambda_2(E_1)$ and $\lambda_2(E_2)$ are positive. Whether there is a unique positive globally stable equilibrium for the robustly permanent system depends in a delicate way on the system parameters, see [38]. For slowly dispersing populations (i.e. $0 < e^j_i \ll 1$ for all $i, j$), the robust permanence condition is particularly

Straight forward to verify. Let $E^j_i$ be the largest solution to $x^j_i \quad f_i^j \ (x^j_i) = 0$ i.e. the equilibrium attained by species $I$ in patch $j$ when there is no dispersal and no competitors. For slowly dispersing populations, $E_1$ is close to the product $(E_1^1, \ldots, E_1^m)$, and the matrix $A^2(E_1)$ is close to the diagonal matrix $\mathrm{diag}(f^i_2(E_1{}^j))^m_{j=1.}$ Hence the invasion rate $\lambda_2(E_1)$ is close to $\max_j f_2^j(E^j_1) > 0$. So we obtain: The system with two slowly dispersing competing species is robustly permanent if and only if system refers to system of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

(8) $\max_j f_2^j(E^j_1) > 0$ and $\max_j f_2^j(E^j_2) > 0$

Intuitively, robust permanence requires for each species there is at least one patch where it can persist.

For *systems, particle interactions, financial transactions, cosmological interactions and reactions, as we shall show model by model the* situation are even easier. Since (under mild assumption on the prey dynamics) the Morse decomposition of $\partial C$ is given by two equilibria $\{0,E_1\}$, the spatial predator-prey system is robustly permanent, if the predator can invade at the prey equilibrium: $\lambda_2\ (E_1) > 0$. However, the global dynamic are likely to be more complicated. For instance, if in each patch there is a globally stable limit circle (e.g.,[39-41]) and dispersal is sufficiently slow (i.e. $0< e^j{}_i \ll 1$for all *i, j*), our result about robust permanence and the theory of normally hyperbolic manifolds [42] imply that there is a positive *m* dimensional torus, which attracts almost all orbits in C, whenever $\max_j f_2^j (E^j_1)>0$

**5.2 Rock-Paper-Scissors Dynamics.** The Lotka-Volterra model of rock-paper-scissor dynamics is a simple model that is used as property for understanding intransitive ecological outcomes [43,44]. Here, a simple spatial version of this dynamic is given by

$$\frac{dx_1^j}{dt} = x_1^j(1-x_1^j - \beta^j\ x_2^j\ - \alpha^j\ x_3^j\ ) + \sum d_1^{jk}\ \ x_1^k\ - e_1^j\ x_1^j$$

$$\frac{dx_2^j}{dt} = x_2^j(1-\alpha^j\ x_1^j\ - x_2^j - \beta^j\ \ x_3^j\ ) + \sum d_2^{jk}\ \ x_2^k\ - e_2^j\ x_2^j$$

$$\frac{dx_3^j}{dt} = x_i^j\,(1-\beta^j\ x_1^j\ - \alpha^j\, x_2^j - \beta^j\ \ x_3^j\ ) + \sum d_3^{jk}\ \ x_3^k\ - e_3^j\ x_3^j$$

Where $a^j \in (0,1)$ and $\beta^j > {}^1$ for all *j*. A more general version of this dynamic is presented in [44]. Under the assumption that there is no cost to dispersal (i.e. $e^j_i = \sum_l d^{l,j}_i$ for all *i, j*) , the maximal invariant set in $\partial C$ consist of the origin 0 and a heteroclinic cycle connecting positive single species equilibria$E_1$, $E_2$ and $E_3$. For slowly dispersing population ($0< e^j_i \ll 1$ for all *i,* j), $E_1$ is close to $(x_1, x_2, x_3) = (1,\ldots,1,0,\ldots,0,0,\ldots,0)$ and the invasion rates $\lambda_2\ (E_1)$and $\lambda_3\ (E_1)$are close to $\max_j 1-a^j$ and $\max_j\ 1 - \beta^j$. For the other equilibria $E_i$, the debitable rates of the missing species are also given by $\max_j 1-a^j$ and $\max_j\ 1 - \beta^j$. Consider the Morse decomposition of $\partial C$ given by 0 and the heteroclinic cycle. Since $\lambda i\ (0) = 1> 0$ for $i = 1, 2, 3, 0$ is robustly insaturated. An algebraic computation reveals that (4) holds for all invariant supported by the heteroclinic cycle (i.e. all convex combination of the Dirac measures supported by the equilibria) if and only if

$$\max(1-a^j) > -\max\ (1 - \beta^j\ )$$

$j$ $j$

Equivalently,

$$2 > \min a^j + \min \beta^j$$

$j$ $j$

In particular, even if the heteroclinic cycle is attracting for each patch when the system is uncoupled (i.e. $2 < a^j + \beta^j$ for all $j$ ), it can be repelling for the weakly coupled system.

5.3. Three species **Lotka Volterra in spatially homogenous environments.** Consider Lotka-Volterra dynamics in a spatially homogenous environment: $(x + a)^n =$

(9) $$\frac{dx_i^j}{dt} = x_i^j (r_i - \sum_s a_{is} x_s^j) + \sum_l d_i^{jl} x_i^l - e_i^j x_i^j , \quad i = 1, \dots, k; \; j = 1, \dots, m$$

We assume that dissipative can be shown by a linear Liapunov function: There are $c_i > 0$ such that $\sum_i c_i a_{is} x_i x_s < 0$ holds for all $x \in R_+^k \setminus \{0\}$. Then the weighted sum of densities across all patches $\sum_{i,j} c_i x_i^j$ is decreasing for large densitities and hence (9) generates a dissipative semiflow with global attractor $r$ . We further assume that for each patch *j* and each species *i* the immigration rate equals the emigration rate , i. e,

(10) $$e_i^j = \sum_l d_i^{jl}$$

This guarantees that the set H= $\{x \in C: x_i^j = x_i^l \forall i, j, l\}$ of spatially homogenous states is forward invariant under (9).

Now consider the Lotka-Volterra dynamics without spatial structure

(11) $$\frac{dx_i}{dt} = x_i(r_i - \sum_s a_{is}x_s), \qquad i = 1, \dots, k$$

(11) $$\frac{dx_i}{dt} = x_i(r_i - a_{is}x_s), \quad i = 1, \dots, k$$

It has been shown [12,ch.16.1,16.2] That for $k$ =3,(11) is robustly permanent (robust meaning here within the class of Lotka –Volterra systems) if and only if all equilibria on the boundary of $R^3_+$ are unsaturated and whenever there is a heteroclinic cycle connecting the one species equilibria then this cycle is repelling. Equivalently the following four conditions hold. There exists an interior equilibrium $\dot{x}$ (i.e $A\dot{x} = r \;\; with \; \dot{x} \gg 0$); Det (-A)>0; the 2 species subsystems are not bistable competition systems if there is a heteroclinic cycle between the one species equilibria, say $\mathbf{E_1 \to E_2 \to E_3 \to E_1}$ then the following inequality holds

12) $$\lambda_2(E_1)\lambda_3(E_2)\lambda_1(E_3) > |\lambda_3(E_1)\lambda_1(E_2)\lambda_2(E_3)|$$

Furthermore, the proof in [12, ch.16.1, 16.2] shows that the boundary flow of (11) has a simple Morse decomposition: if (iv) applies then there are two Morse sets, the origin 0 is a repeller, and the heteroclinic cycle is the dual attractor (within$\partial R^3_+$); if (iv) does not apply then the finitely many boundary equilibria form a Morse decomposition of the boundary flow.

**Species, population we mean types of and population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.**

**Theorem 4:** Under the above assumptions,(9) for k=3 is permanent if and only if *(i),(ii),(iii)* and *(iv)* hold.

Proof: Suppose (9) is robustly permanent. Since the nonspatial system (11) is the restriction of (9) to the invariant subspace **H**, (11) must be robustly permanent and hence*(i)-(iv)* hold

Conversely, suppose that (i)-(iv) hold. We show that the above Morse decomposition of the boundary flow of (11) (now in **H**) is also Morse decomposition for (9) restricted to the maximal invariant subset of $\partial C$ which is contained in the spatial two species subsystems of (9). For this we show that in each 2 species system the orbits of (9) converge to **H** (For the one species subsystems this is obvious.) If a two species subsystem of (11) has an internal equilibrium, say $\mathbf{E_{12}}$ then by assumption (ii),this is not a saddle but a sink, and hence $a_{11}a_{22} > a_{12}a_{21}$. This implies that the 1-2 submatrix of A is VL-stable [12,ch.15.3,],and [45] implies that $\mathbf{E_{12}}$ is the global attractor for the spatially structured system (9) restricted to the first two species.

If a two species subsystem of (11) has no internal equilibrium, then a one species equilibrium say $\mathbf{E_1}$ is the global attractor. Then Lemma 2 (see below) show the corresponding result for (9).A similar proof applies to the case where the origin **0** is the global attractor of a two species subsystem.

Finally observe that under the above assumptions (9)-(10), the invasion rates $\lambda_i$ at spatially homogenous boundary equilibria are the same for (9) and (11) Since all boundary equilibria are unsaturated for (11), they are for (9) as well. The same applies in case (iv) to the heteroclinic cycle. Note that the condition (12) on the eigenvalues is equivalent to the fact that all invariant measures supported on the heteroclinic cycle are unsaturated, see [16,Ex.4.5]

Species we mean types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

**Lemma 2:** Suppose, in a two species Lotka –Volterra system (11) (k=2) all interior orbits converge to the one species equilibrium $\mathbf{E_1}$ .Then all orbits of the spatial version (9) (k=2) converge to the spatially homogeneous one species equibrium $\mathbf{E_1} \in \mathbf{H}$.

**Proof:** If the system is of predator-prey or mutualistic type then Hasting's [45] result applies again. So let us assume that the local interaction is given by the competition system

(13) $\frac{dx_1}{dt} = r_1 x_1(1 - x_1 - ax_2)$ $\frac{dx_2}{dt} = r_2 x_2(1 - \beta x_1 - x_2)$

With 0<a <1< $beta$.In this case, Kazuo Kishimoto (letter to JH, Oct 1987) has given the following argument. There exists a family of forward invariant rectangles contracting to $\mathbf{E_1}$. This shows that all interior solutions of (9) (k=2) converge to $\mathbf{E_1}$.

Theorem 4 shows that a spatial network of identical patches is permanent, if the within patch dynamics is permanent, and three species Lotka-Volterra dynamics. It is not clear how to extend this result to more than three species. The crucial step in the three species case is that permanence precludes bistable two species subsystems. The spatial version of a bistable two species system allows plenty of stable equilibria outside **H,** see [33].However, permanent four species Lotka Volterra systems may have two species bistable subsystems [12,ch.16.4].So there may be invariant sets on the boundary outside **H** which need to be unsaturated to make (9) permanent.

## 6. MORE APPLICATIONS

6.1. **Disease, particle, product, marketing, financial, econometric dynamics with density-dependent demography**. GAO and Hethcote [46] introduced a model of disease dynamics with density-dependent demography. Here we consider a variation of their model in which a population of size P has a per-capita birth rate b(P) and per –capita death rate d. To ensure that the population persists in the absence of the disease, we assume that there exists K>0 such that b(K)=d. To describe the disease dynamics, let *S* be the number of individuals susceptible to the disease, *I* the Number of infected individuals, and R the number of individuals that have recovered from the disease. We assume that P=S+*I* +R. On the same lines we can think of sick and non performance accounts, and those which are conducting themselves well, and those which are satisfactory and find a parallel formula. The same thing holds true in case of a black hole eating matter and black hole radiating radiation.. If the disease transmission is asymptotic, β is the contact rate between susceptible and infected individuals, γ is the rate at which individuals recover from the disease, and m the mortality rate due to disease, then the population dynamics is given by

$$\frac{dp}{dt} = b(P)P - dP - mI$$

$$\frac{dI}{dt} = \beta(P - I - R)\frac{I}{\epsilon + P} - (\gamma + d + m)I$$

$$\frac{dR}{dt} = \gamma I - dR$$

Where ϵ>0 is a constant It is useful to introduce a change of coordinates in which $y=\frac{I}{P}$ and $z=\frac{R}{P}$:

$$\frac{dP}{dt} = (b(P)-d-my)P$$

$$\frac{dy}{dt} = \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(P)\ y$$

$$\frac{dz}{dt} = \gamma y - (b(P) - my)z$$

Setting $x_1$=P and $x_2$=(y,z) 'where' denotes transpose yields a structured model where $A_1(x)$=b-d-my and

$$A_2(x) = \begin{pmatrix} \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(p) & O \\ \gamma & my-b(p) \end{pmatrix}$$

Since {0,(*K*,0,0)}is a Morse decomposition of the boundary dynamics Theorem 3 and Remark 1 imply that this model is robustly permanent if and only if b(0)>d and $\beta\frac{K}{\epsilon+K} > \gamma + m + d$. In particular, for $\epsilon > 0$ sufficiently small, one requires that the basic reproductive number $\frac{\beta}{r+m+d}$ is greater then one.

Remarkably this same criterion determines robust permanence of model with significantly more complicated boundary dynamics. For example, suppose the focal population is a predator species. If the prey has abundance *N* and exhibits logistic dynamics f(*N*)is the per-capita predator consumption rate of the prey and b(N)is the per-capita reproductive rate of the predator then the dynamics become

$$\frac{dN}{dt} = rN(1 - N/K - f(N)P$$

$$\frac{dP}{dt} = (b(N) - d - my)P$$

$$\frac{dy}{dt} = \left( \beta(1 - y - z)\frac{P}{\epsilon+P} - r - m(1 - y) - b(N)\ y \right)$$

$$\frac{dz}{dt} = \gamma y - (b(N) - my)z$$

Where r>0 and *K*>0, and *f*(*N*),b(N)satisfy f(0)=b(0)Define $x_1 = N$, $x_2 = p$, $x_3 = (y,z)'$ Let us assume that the disease-free community (i.e.y=z=0)is permanent (i.e.b(K)>d) and let A be the global attractor in the interior of the P-N plane Under these assumption a Morse decomposition of the boundary dynamics is given by $(x_1, x_2, x_3) = (0,0,0)$, $(x_1, x_2, x_3) = (K, 0,0)$and A. Since the equilibria (0,0,0) and (k,0,0) are unsaturated, it remains to characterize $\lambda_3(\mu)$for any invariant measure μ supported by A. Proposition 3 and Remark 1 imply

$$\lambda_2(\mu) = \int_c (bx_1)\text{-}d)d\mu(x) = 0$$

*Hence* $= \int_c bx_1)d\mu(x) = d$ *and*

$$\lambda_3(\mu) = \beta \int \frac{x_2}{\epsilon+x_2} d\mu(x) - r - m - d$$

**Theorem 3**: and Remark 1 implies that this structured model is robustly permanent if $\frac{\beta}{r+m+d} > 1 and\ \epsilon > 0$ is sufficiently small.

6.2. Gene networks. The repressilator is an oscillatory gene network based on three (or more generally an odd number of transcriptional repressors. Such a system has been genetically engineered in E-coli by Elowitz and Leibler [47] Mathematical model go back to [48.49] they all involve concentrations of proteins and mRNs. A modified model with auto-activation was suggested in [50]

$$(14)\ x_i = \beta(y_i - x_i)$$

$$(15)\ y_i = aF(x_{)i} - y_i)$$

With

$$(16)\quad F(x)_i = x_i g\ (x_i, x_i - 1) = \frac{x_i}{1+x_i+p x_{i-1+Kp\, x_i x_i - 1}}$$

Here $x_i$ and $y_i$ are normalized concentrations of proteins and mRN As belonging to gene $i$,p is the strength of repression and k a cooperativity parameter this is a structured system of type(2)

$$(17)\ \begin{pmatrix} x_i \\ y_i\ dot \end{pmatrix} = \begin{pmatrix} -\beta & \beta \\ ag\,(x_i, x_i) & -1 \end{pmatrix} \begin{pmatrix} x_i \\ y_i \end{pmatrix}$$

In [50,Thm 3] it was shown that for n=3,a>1 and p>1,system(14)has a heteroclinic cycle connecting the 3 single gene equilibria (similar to the rock –scissors-paper dynamics of section5.2)Moreover

- the system is permanent i.e. the boundary of $R^6$+is repelling, if $\lambda + \mu > 0$
- this heteroclinic cycle is asymptotically stable if $\lambda + \mu < 0$

where $\lambda > 0$ and $\mu < 0$ are the invasion rates of the two missing genes i.e. the leading eignvalues of the 2X2 matrix in (17)evaluated at the appropriate single gene equilibria (for explicit expressions see[50,eqs.(127,128)].)

Theorem 3 implies the permanence condition implies the stronger conclusion of robust permanence A similar result holds for n odd, but the heteroclinic cycle is then actually a heteroclinic network.

Note: For references please see the original paper.

One of the examples for which we apply the model are as follows: **Progress in Particle and Nuclear Physics Volume 88, May 2016, Pages 1–28 Chiral magnetic and vortical effects in high-energy nuclear collisions—A status report D.E. Kharzeeva, b, J. Liaoc, d, , , S.A. Voloshine, G. Wangf**

The interplay of quantum anomalies with magnetic field and vorticity results in a variety of novel non-dissipative transport phenomena in systems with chiral fermions, including the quark–gluon plasma. Among them is the Chiral Magnetic Effect (CME)—the generation of electric current along an external magnetic field induced by chirality imbalance. Because the chirality imbalance is related to the global topology of gauge fields, the CME current is topologically protected and hence non-dissipative even in the presence of strong interactions. As a result, the CME and related quantum phenomena affect the hydrodynamical and transport behavior of strongly coupled quark–gluon plasma, and can be studied in relativistic heavy ion collisions where strong magnetic fields are created by the colliding ions. Evidence for the CME and related phenomena has been reported by the STAR Collaboration at Relativistic Heavy Ion Collider at BNL, and by the ALICE Collaboration at the Large Hadron Collider at CERN. The goal of the present review is to provide an elementary introduction into the physics of anomalous chiral effects, to describe the current status of experimental studies in heavy ion physics, and to outline the future work, both in experiment and theory, needed to eliminate the existing uncertainties in the interpretation of the data. Keywords Chiral magnetic effect; Chiral vortical effect; Chiral anomaly; Quark-gluon plasma; Heavy ion collisions Corresponding author at: Physics Department and Center for Exploration of Energy and Matter, Indiana University, 727 E Third Street, Bloomington, IN 47405, USA. Copyright © 2016 Elsevier B.V. All rights reserved.

## Sufficient Conditions For The Existence And Linear Stability Of Time Independent Age Distributions In Respect Of Particles, Propositions, Problems, Hegelian Structures, Deleuzean Treatises, Interactions And Transactions In Cosmos And Others

**Abstract:** Gopalaswamy studied sufficient conditions are obtained for the existence and linear stability of time independent age distributions in two species competition with age time lagged density dependent mortality and fertility functions.. Work was published in Bulletin Australian Mathematical Society Volume 25 (1982). We show how the model could be applied in case of finance, econometrics, psycho economics, interactions between particles and transactions in nature. Basically the competition is for resources, notwithstanding such experiments could be conducted in laboratory like a rat kept in front of a cat, such incidents could be represented mathematically, and the corresponding linear stability conditions ca be obtained in case of interactions in cosmos. Take for example the case of a cash credit account of a borrower. Many inward bills waiting for paid, or inland or foreign bills payment all are in line for the debit of the account, which also acts as a resource. In addition there are recurrent cases where in the credit is issued to the cash credit account in case of demand and inland bills. In case of economics, investment has inventories and it has to be utilized towards the end of consummation of output. And this is case of output gobbling part of investment which is inventories. Output has to be sold to the customer and the amount received is diverted partly to savings and consumption and part could be employed towards the consolidation of the system. So this is a case of investment itself acting as a both receiver and payer. It shall eat and also produce the finance. Savings as stated above could be deployed for the further diversification and proliferation and that is towards the end of putting the company on terra firma. Another example in cosmology is the case of black holes guzzling matter. There are lot

of black holes and lot of matter. Result of such a guzzling the production of Hawking radiation. Horava gravity is another example. How the conformal field theory could be changed in the case of Horava gravity. Still another case is the case of Vonneumann algebra in the cae of Horava gravity. In this case like cash credit account, the Horava gravity and Von Neumann algebra act as both receivers and takers. Such incidents are legion in cosmology and even in philosophy. We have considered self as individual general ledger, and the debits and credits are raised to the same account thus anointing as both a receiver and payer. We give plethora of such cases and apply the model. Due to spatial constants, same module is used for all the modules and we rewrite to state that the model is applicable to all modules. This is due to spatial restrictions. What is to be done is that each and every parameter is to be changed with a subscript or a superscript and finally it is to be concatenated towards the end of obtention of prediction model or optimizing model, for Duality of motion model for further analysis. Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would

have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

**Introduction**

Let $P_1$ *(t)* and $P_2$ *(t)* denote the total particles, interactions, transactions, etc, which we dilate upon further investigation. Perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. **Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems** sizes (or biomasses) at time $t \geq 0$ of two interacting species living in a common habitat and competing for a common pool of resources. The competition for resources will be implicit in our model similar to that in the two species Lotka- Volterra competition system. Assuming constant sex ratios in the two species we can consider $P_1$ *(t)* and $P_2$ *(t)*to be the Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

population of f; immigration, emigration and internal dispersion in the habitat are assumed to play no significant role in the dynamics of the community.Popoulation we mean Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

We suppose that the two species contain respectively $p_1(a, t)\,da$ and $p_2(a, t)\,da$ individuals with ages between $a$ and $a + da$ $(a \geq 0)$ at time $t$ so that we have

(1.4) $P_1(t) = \int_0^\infty \rho_1(a, t)\,da,\ P_2(t) = \int_0^\infty \rho_2(a, t)\,da,\ t \geq 0.$

The rates of change of the two species densities are by definition $D\rho_1(a, t)$ and $D\rho_2(a, t)$ where

(1.5) $D\rho_i(a, t) = \limsup_{h \to 0} \frac{\rho_i(a+h, t+h) - \rho_i(a, t)}{h},\ i = 1, 2.$

Assuming the existence of the limits in (1.2) we consider the following time lagged model system

(1.6) $$\begin{cases} D\rho_1(a,t) + f_1(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_1(a, t) = 0, & a>0,\ t>0 \\ D\rho_2(a,t) + f_2(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_2(a, t) = 0, & \text{(1.4)} \\ \rho_1(0, t) = \int_0^\infty b_1(a, P_1(t), P_2(t-\tau))\rho_2(a, t)da, & t>0, \\ \rho_2(0, t) = \int_0^\infty b_2(a, P_1(t), P_2(t-\tau))\rho_1(a, t)da, & \end{cases}$$

(1.5) $\rho_i(a, s) = \varphi_i(a, s),\ a.0,\ s \in [-\tau, 0],\quad i = 1, 2,$

(1.6) $P_i(s) = \int \varphi_i(a, s)\,da,\ s \in [-\tau, 0],$

Where $\tau$some is fixed nonnegative constant; $f_1$ and $f_2$ denote the age and density dependent mortality rates with time lags in interspecific interactions, $b_1$ and $b_2$ denote the age and density dependent fertility functions again having time lags in the interspecific interactions; $\varphi_1$ and $\varphi_2$ denote the initial age distributions needed for the formulation of the model. Note we mean the population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

The model system (1.1)-(1.6) will be meaningful only if $\varphi_1, \varphi_2, f_1, f_2, b_1, b2$ are nonnegative functions of the respective arguments. Also since the total initial populations have o be finite, $\varphi_1(\cdot, s)$ should belong to $L_1(R^+)$ for each $s \in [-\tau, 0]$. That $f_i$ and $b_i$ $(i = 1, 2)$ are independent of t explicitly and this assumption can be interpreted to represent the temporal uniformity of the environment.

One of the fundamental questions for (1.1)-(1.6) is the following; under what conditions on $f_i, b_i$ and $\tau$ does the system (1.1)-(1.4) have time independent nontrivial solutions, and when such solutions, and when such solutions exist, are they stable with respect to some suitable stability criterion? When such stable time invariant solutions exist for (1.1)-(1.4) we say that the two types has a stable age distribution. Existence of a stable age distribution is the analogue of the existence of a stable steady state for the corresponding age independent system. Types we mean

types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

In this article we investigate the above question for the case of a competitive interaction with some additional constraints on the vital rates $f_i$ and $b_i$ $(i = 1, 2)$. Age dependent population systems without time delays have been considered by several authors (Gurt in and MacCamy [3, 4, 5], Gurt in and Levine [6], Haimovici [7] and Rotenberg [8]). The question of the relation between an age independent system and a corresponding age dependent system has been considered by Gurt in and MacCamy [4] and Gurt in and Levine [6] who have established an asymptotic (*as* $t \to \infty$ ) relation between such models by constructing a higher dimensional 1 umped parameter system (in terms of ordinary differential equations) to represent the age dependent distributed parameter system. Assuming the existence of stationary age distributions Haimovici [7] considers their stability in a system of two interacting populations explicitly taking into consideration the dynamical nature of the encironments resources and pollution.

In this article we add another realistic feature namely time lags in the age dependent models and obtain sufficient conditions for the existence of stationary age distributions in (1.1)-(1.4) and show that the time lags in interspecific interactions have no effect on the linear stability of the age distributions although the decay rates of perturbations will depend on the magnitudes of the time delays. Such "harmless" natures of time lags in interspecific interactions have been noted by the author for the age independent models in [1, 2].

## 2. Existence and uniqueness of solutions

We first narrate our assumptions on the vital rates of $f_i$ and $b_i$ $(i = 1, 2)$ so that the system (1.1)-(1.6) will denote a competition system. Let $BC\ (0, \infty)$ denote the linear space of bounded continuous real functions on $(0, \infty)$ with the norm defined by

$$||f|| = \{max\ |\ f(t)\ |,\ t \in (0, \infty)\}.$$

Let $K_1$ be the cone of nonnegative functions of $BC\ [0, \infty]$. Let $C\ [-\tau, 0]$ denote the space of real continuous functions with a norm defined by $||g|| = \{\ max\ |g(s)\ |,\ s \in [-\tau, 0]\}$; let $K_2$ be the nonnegative cone of functions in $C\ [-\tau, 0]$ . We present our first set of assumptions on $f_i$, $b_i$ and $\varphi_i$ for $i = 1, 2$.

$A_1$. $f_1 \in C\ (R^+ \times K_1 \times K_2, R^+)$, $f_1 \geq d_1^*$ on $R^+ \times K_1 \times K_2$

$f_2 \in C(R^+ \times K_2 \times K_1, R^+)$, $f_2 \geq d_2^*$ on $R^+ \times K_2 \times K_1$,

( $d_1^*$, $d_2^*$ are positive constants);

$f_1$ and $f_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ ; that is, *for* $i = 1, 2$,

$$|f_i(a, P_1', P_2') - f_i(a, P_1'', P_2'')| \leq d_i [||P_1' - P_1'', P_2' - P_2''||]$$

($d_1$, $d_2$ are positive constants),

$$\partial f_1/\partial P_1, \partial f_1/\partial P_2 \in C(R^+ \times K_1 \times K_2, R^+),$$

$$\partial f_2/\partial P_1, \partial f_2/\partial P_2 \in C(R^+ \times K_2 \times K_1, R^+),$$

$A_2$. $b_1 \in C(R^+ K_1 \times K_2, R^+)$, $b_1 \leq \beta_1^*$ *on* $R^+ \times K_1 \times K_2$,

$b_2 \in C(R^+ K_2 \times K_1, R^+)$, $b_2 \leq \beta_2^*$ *on* $R^+ \times K_2 \times K_1$,

*($\beta_1^*$, $\beta_2^*$ are positive constants);*

$b_1$, $b_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ *; that is for* $i = 1, 2$,

$|b_i(a, P_1', P_2') - b_i(a, P_1'', P_2'')| \leq \beta_i [||P_1' - P_1''|| + ||P_2' - P_2''||]$

$B_1$, $B_2$ Are positive constants),

$b_i(\circ, (P_1, P_2)) \epsilon L^{\infty}(R^+, R^+)$ For fixed $P_1, P_2, i = 1,2$,

$$\frac{\partial b_1}{\partial P_2}, \frac{\partial b_1}{\partial P_2} \epsilon\ C\ R^+ \text{x}\ k_1 \text{x}\ k_2, \text{R}^-),$$

$$\frac{\partial b_2}{\partial P_1}, \frac{\partial b_2}{\partial P_2} \epsilon\ C\ R^+ \text{x}\ k_2 \text{x}\ k_1, \text{R}^-),$$

$A_3$. $\varphi_i(\cdot, s) \in L^2(R^+, R^+) \cap C(R^+, R^+)$ for each $s \in [-\tau, 0]$,

$i, = 1, 2,$

By a solution of $(1.1) - (1.6)$ on $[0, T]$, we mean

$P_i: R^+ \times [-\tau, T] \to R^+$ $(i, = 1, 2,)$ With the following properties:

$S_1$. $P_i(\cdot, t) \in L^1(R^+, R^+)$ for each $t \in [-\tau, T]$;

$S_2$. $P_i(a, \cdot) \in C([-\tau, T], R^+)$ for each $a \in R^+$;

$S_3$. $DP_i$ Exists along the characteristics $t - a =$ constant

On $R^+ \{R^+ \cap [0,T]\}$ And is continuous $(i, = 1, 2,)$;

$S_4$. $P_i$ Satisfies the system (1.1)- (1.6) for

$(a,t,) \in R^+ [-\tau, T]$ .

Let us first convert the system (1.1)- (1.6) into an equivalent system of integral equations. For brevity we denote by P the pair ( $P_1, P_2$) If we let

$$P_i(a+x, t+x) = P_i(x), (i, = 1, 2,),$$

Then (1.3) considered at $(1.3)$ considered at $(a+x, t+x)$ becomes

$$\frac{dp_i(x)}{dx} + f_i(a+x, P(t+x)) \; P_i(x) = 0$$

Which has a solution of the form?

$$P_i(x) = P_i(0) \exp\left[ - \int_0^\infty f_i(a+s, P(t+S)) ds \right]$$

And hence

(2.1) $p_1(a+x, t+x) = p_1(a,t) \, exp\left[ - \int_o^x f_1(a+s, p(t+s)) ds \right]$

For $a \geq t$ *we choose (a,t)=(a-t,0)* in (1.3) and $x=t$ in (2.1) so that

(2.2) $p_1(a,t) = \varphi_1(a-t, 0) \exp\left[ - \int_o^t f_1(a-t+s, p_1(s)) ds \right]$

$$O \leq t \leq a, i = 1,2,$$

With

$$P_1(s) = (p_1(s), p_2(s-t)),$$

$$P_2(s) = (p_1(s-t), p_2(s)),$$

For $0 \leq a < t$ we choose *(a,t)=(0,t-a)in (1.3)and x=a in (2.1)* so that

(2.3) $p_i(a,t) = B_i(t-a) exp[-\int_o^a f_1(s, p_i(t-a+s)) ds]$

$$t > a \geq, i = 1,2,$$

Where

| |
|---|
| (2.4) $B_i(t) = p_i(0, t), i = 1,2, t > 0.$ |
| Define $M_i and\ L_i\ (i = 1, 2)$ as follows: |
| $M_i\ (a,t, p_i(t)) = exp[-\int_o^{t-a} f_1(s, p_i\ (a - a + s))ds]$ |
| (2.5) $L_i(a, t,\ p_i(t)) = \exp[-\int_o^t f_1(a+s,\ p_i\ (s))ds]$ |
| It will now follow from (1.1),(1.3),(2.2)-2.5)that |
| (2.6) $p_i(t) = \int_o^t B_1(a)\, M_1(a, t, p_i\ (t)da + \int_o^\infty \varphi_i\ (a, 0)\ L_i(a, t, p_i(t)da,$ |
| $t > 0, i = 1,2,$ |
| (2.7) $p_i(s) = \Phi_i\ (s) = \int_o^\infty \varphi_i\ (a, s)da,\ s \epsilon [-\tau, 0], i = 1,2,$ |
| (2.8) $\beta_i$(t) = $\int_0^t \beta_i$(a)$b_i$(t-a, $P_i(t)$)$M_i$. $\left(a, t, P_i(t)\right)da$ |
| $+\int_0^\infty b_i\left(a + t, P_i(t)\right)\varphi_i(a, 0)L_i\left(a, t, P_i\ (t)\right)da$ |
| $i$ = 1, 2, t > $0$ |
| the equivalence of (2.6)-(2.8) with (1.1)-(1.6) is established by the following whose proof is identical to a similar result of Gurtin and MacCamy [4]; hence we omit the details of the proof of the following. |
| THEOREM 1 : *Suppose the system* (1.1)-(1.6) *satisfies the conditions* $A_1, A_2, A_3$ . *If* $(P_1, P_2)$ *is a solution of* |
| *(1.1)-(1.6) then the total population* $P_1, P_2$ *and the birth rates* $B_1, B_2$ *satisfy the integral equation (2.6)-(2.8).* |
| *conversely if* $P_1, P_2, B_1, B_2$ *are nonnegative continuous solution of (2.6)-(2.8) on [ 0, T] and if* $P_1, P_2$ *are* |
| *defined by (2.2)-(2.4) on* $R^+$ x [0, T] *then such* $P_1, P_2$ *provide a solution of (1.1)-(1.6) on* $R^+$ x [0, T] . |
| The following a *priori* estimates are useful to prove our existence theorem below. It will immediately follow from the bounds for $f_i$ and $b_i$ that |
| |
| $M_i(a, t, P) \leq \exp[-d_i^*\ (t - a)]\ \leq\ 1, if\ \ t \geq a,$ |
| *(2.9)* |
| $L_i(a, t, P) \leq \exp[-d_i^* t]$ |
| From (2.7)-(2.8) we derive that |

$$\beta_i(t) \le \int_o^t \beta_i^* \beta_i(a) \exp[-d_i^*(t-a)]\, da + \beta_i^* \Phi_i(0) \exp[-d_i^* t], \quad i = 1,2,$$

And hence by Gronwall's inequality,

(2.10) $$\beta_i(t) \le \beta_i^* \Phi_i(0) \exp[\delta_i t], \ \delta_i = \beta_i^* - d_i^*$$

(2.5), (2.6) and (2.10) lead to

(2.11) $$|P_i(t) - 0_i(0)| \le \beta(0)\, exp[\delta_i\, a - d_i^*(t-a)]\, da$$

$$+\int_0^\infty \varphi_i\ (a,0)|1 - exp(-d_i^* t)|da$$

$$\le \beta_i^* \Phi_i\ (0)\, exp(\beta_i^* t) \left\{\frac{1-exp[-(d_i^*+\beta_i^*)t]}{d_i^*+\beta_i^*}\right\} + \beta(0) d_i^* t$$

$$\le [\beta_i^* \beta(0)\, exp(\beta_i^* t) \beta(0) d_i^*] t \quad for\ t > 0.$$

We can now prove the following result on the existence and uniqueness of solution of the integral equations (2.)-(2.8)

**THEOREM 2:** *Assume that the system* (1.1)-(1.6) *satisfies the hypotheses*

$A_1, A_2, A_3$. *Then there exists a unique set of continuous function* $(P_1, P_2, B_1, B_2)$ such *that*

$$P_i : R^+ \to R^+, \qquad B_i : R^+ \to R^+ \ (i = 1,2),$$

*Which satisfy the integral equations* (2.6)-(2.8)

Proof. We will first show the existence of local solution on same (possibly small) interval [0, T] $\subset R^+$ such that

$$P_i \in C([0,T], R^+), \quad B_i \in C([0,T], R^+), i = 1,2,$$

Satisfying (2.6)-(2.8) and then show the $P_i$ and $B_i$ can be continued as solution for any finite interval in $R^+$ for some fixed

Positive number $T$ we let

$$C^+[-\tau, T] = \{f \in C([-\tau, T], R) | f \ge 0\}$$

Let $||.||_T$ denote the supremum norm in $C^+[-\tau, T]$. Then for any

$(x, y) \in C^+[-\tau, T] \times C^+[-\tau, T]$ we define

$||(x,y)|| = ||x||_T + ||y||_T$,

$C_T = C^+[-\tau, T]$, $C_T^2 = C^+[-\tau, T] \times C^+[-\tau, T]$.

Members of $C_T$ will be denoted by $Z_t$ with the meaning that

$Z(t+s) = z_t(s)$ *for* $s \in [-\tau, T]$ *and* $t \in [0, T]$. Now for any fixed

$P_t = (P_{1t}, P_{2t}) \in C_T^2$, (2.8) is a linear system of uncoupled Volterra integral equations in B = $(B_1, B_2)$ and hence (2.8) will

have a unique continuous solution under the assumption $A_1 - A_3$. let us denote such a solution by

$B(t) = B^T(P_t), \quad t \in [0, T]$

Or

$B_1(t) = B_1^T(P_{1t}, P_{2t}), \quad B_2(t) = B_1^T(P_{1t}, P_{2t}), \quad t \in [0, T]$

If we supply this solution $(B_1, B_2)$ in (2.6) we find that (2.) will be satisfied if and only if $(P_{1t}, P_{2t})$ is a fixed point of the

operator

$\pi(P_t) = \{\pi_1(P_{1t}, P_{2t}), \pi_2(P_{1t}, P_{2t})\}$

Where $\pi_1$ *and* $\pi_2$ respectively denote the right sides of (2.6) for *i =1,2* .

It is not difficult to see from the nature of $M_i$, $L_i$ and the hypotheses on $f_i, b_i, \varphi_i$ that

$B: C_T^2 \to C_T^2$ *and* $\pi: C_T^2 \to C_T^2$

(which guarantees the nonnegativity of $P_i$ and $B_i$). We will first show that

There exist positive constants $r^*, t^*$ such that $\pi$ is a mapping of a sphere

$$S_{t*}(\Phi, r^*) = \left\{ P_t = (P_{1t}, P_{2t}) \in C_{t*}^2 \mid \|P_t - \Phi_t\|_{t*} \le r^* \right\}$$

Into itself and also is a contraction where $\Phi_t$ is defined on [$-\tau$, $t^*$]

By the following:

$\Phi_t = (\Phi_{1t}, \Phi_{2t})$;

(2.12)

$$\Phi_{it} = \begin{cases} \int_0^\infty \varphi_i(a,s)\,da \ for\ s \in [-\tau, 0]\,. \\ \int_0^\infty \varphi_i(a,0)\,da \ for\ s \in [0, t^*]\,. \end{cases}$$

We note that similar to (2.11) one gets the estimates

$$|\mathcal{P}(P)(t) - \Phi_i(0)| \le \{B_i^*\Phi_i(0)\exp[B_i^* t] + \Phi_i(0)d_i^*\}t$$

Which together with (2.7) and (2.12) lead to

(2.13) $$\|\mathcal{P}(P_t) - \Phi_t\|_{t*} \le \begin{cases} \sum_{j=1}^2 \|\Phi_{jt}\|_{t*}[B_j^*\exp(B_j^* t) + d_i^*] \\ t^* \end{cases}$$

Thus if we choose,

$$r^* = T\sum_{j=1}^2 \|\Phi_{jt}\|_{t*}[B_j^*\exp(B_j^* T) + d_i^*]$$

Then $\mathcal{P}$ is a mapping of $S_{t*}(\Phi, r^*)$ into itself for all $t^* \le T$. we then will have to show that for small $t^* > 0$, $\mathcal{P}$ is a contraction.

Let $P_t^{(1)}$ *and* $P_t^{(2)}$ be arbitrarily chosen in $S_{t*}(\Phi, r^*)$ and consider $|\mathcal{P}(P_t^{(1)}) - \mathcal{P}(P_t^{(2)})|$ for t $\le t^*$.

Corresponding to the chosen $P_t^{(1)}$, $P_t^{(2)}$ we let $B^{(1)}(a) = B(P_T^{(1)})$ and $B^{(2)}(a) = B(P_T^{(2)})$. Then from

(2.6),

$$|\mathcal{P}(P^{(1)}(t)) - \mathcal{P}(P^{(2)}(t)| \le \left| \int_0^t B_1^{(1)}(a)\left\{M_1\left(a,t,P_t^{(1)}\right) - M_1\left(a,t,P_t^{(2)}\right)\right\}da \right|$$

$$+ \left| \int_0^t [B_1^{(1)}(a) - B_2^{(2)}(a)]\, M_1(a,t,P_t^{(2)})da \right|$$

$$+ \left| \int_0^t \varphi_1(a,0)\left[L_1\left(a,t,P_t^{(1)}\right) - L_1\left(a,t,P_t^{(2)}\right)\right]da \right|$$

$$= J_1 + J_2 + J_3 \text{ (say)}$$

using the elementary inequality

$$|e^x - e^y| \le |x - y| \ for\ x, y \le 0$$

We derive from (2.5) that

$$\left| M_1(a,t,\ p_t^{(1)}) - M_1(a,t,\ p_t^{(2)}) \right|$$

$$\leq \int_o^{t-a} \left| f_1(s,\ p_t^{(1)}(a+s)) - f_1(s,\ p_t^{(2)}(a+s)) \right| ds$$

$$\leq d_1 \| p_t^{(1)} - p_t^{(2)} \|_{t^*} \text{(using } A_1\text{)}$$

*Now,using (2.10),*

$$\left| J_1 \right| \leq \int_o^a \beta_1^* \Phi_1(0)\, e^{\delta 1^a} \left( \int_o^{t-a} d_1 \| P_t^{(1)} - P_t^{(2)} \| ds \right) da$$

$$\leq d_1\, \beta_1^*\, \Phi_1(0)\, t^2 \| P_t^{(1)} - P_t^{(2)} \| e^{\delta 1^t/2}$$

In a similar way

$$\left| J_3 \right| \leq \int_o^{\infty} \varphi(a,o) \left[ \int_o^t \left| f_1(a+s, P_t^{(1)}(s) - f_1(a+s, P_t^{(2)}(s)) \right| ds \right] da$$

$$\leq d_1 \Phi_1(0) \| P_t^{(1)} - P_t^{(2)} \| t$$

Now we have to estimate $\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right|$ in order to estimate $J_2$, from (2.8),

$$B_1^{(1)}(a) - B_1^{(2)}(a)$$

$$= \int_o^a \left[ B_1^{(1)}(\sigma) - B_1^{(2)}(\sigma) b_1 \left( a-\sigma, P_t^{(1)}(a) \right) M_1(\sigma, a, P_t^{(1)}(a) d\sigma + \int_o^a \left[ B_1^{(2)}(\sigma) \right. \right.$$

$$\text{X} \left[ b_1(a-\sigma, P_t^{(1)}(a) M_1(\sigma, a, P_t^{(1)}(a) - b_1 \left( a-\sigma, P_t^{(1)}(a) M_1 \left( \sigma, a, P_t^{(2)}(a) \right) \right] d\sigma \right.$$

$$+ \int_o^{\infty} \varphi(\sigma, o) \left[ b_1 \left( \sigma + a, P_t^{(1)}(a) \right) - b_1 \left( \sigma + a P_t^{(2)}(a) \right) \right], L_1 \left( \sigma, a, P_t^{(1)}(a) \right) d\sigma$$

$$+ \int_o^{\infty} \varphi(\sigma, o) b_1 \left( \sigma + a, P_t^{(2)}(a) \right) \left[ L_1 \left( \sigma, a, P_t^{(1)}(a) \right) - L_1 \left( \sigma, a, P_t^{(2)}(a) \right) \right] d\sigma$$

$$= R_1 + R_2 + R_3 + R_4 \text{ (say)}.$$

By the uniform Lipschitz continuity of $b_1$ in $A_2$ we get

$$\left| b_1 \left( a+\sigma, P_t^{(1)}(a) \right) - b_1 \left( \sigma + a, P_t^{(2)}(a) \right) \right| \leq B_1 \| P_t^{(1)} - P_t^{(2)} \| \text{ for } a\ \varepsilon [o, t^*]$$

And hence we have

$$|R_3(a)+R_4(a)|\leq\beta_1\Phi_1(0)||P_t^{(1)}-P_t^{(2)}||_t^{*}+\Phi_1(0)\beta_1 d_1||P_t^{(1)}-P_t^{(2)}||_t^{*a},$$

$$|R_2(a)|=[\int_0^a[B_1^{(2)}(\sigma)\{b_1(a-\sigma, P_t^{(1)}(a))-b_1(a-\sigma, P_t^{(2)}(a))\}M_1(\sigma, a, P_t^{(1)}(a))$$

$$+b_1(a,-\sigma, P_t^{(2)}(a))\{M_1(\sigma, a, P_t^{(1)}(a))-M_1(a-\sigma, P_t^{(2)}(a))\}]d\sigma]$$

$$+b_1\left(a-0, P_t^{(2)}(a)\right)\{M_1(\sigma, a, P_t^{(1)}(a))-M_1(a-\sigma, P_t^{(2)}(a))\}]d\sigma]$$

$$\leq \int_0^a [\beta_1^{*}\Phi_1(0)\, e^{\delta 1^a}\beta_1\| P_t^{(1)}-P_t^{(2)}\|_{t^*}+\beta_1 d_1\| P_t^{(1)}-P_t^{(2)}\|_{t^*}(a-o)]do$$

$$\leq \beta_1^{*}\Phi_1(0)\, e^{\delta 1^a}\| P_t^{(1)}-P_t^{(2)}\|_{t^*}\{\beta_1 a+d_1 a^2/2\}$$

Thus

$$\left|B_1^{(1)}(a)-B_1^{(2)}(a)\right|$$

$$\leq \int_o^a \left| B_1^{(1)}(\sigma)-B_1^{(2)}(\sigma)\right| B_1^{*d}\sigma$$

$$+\Phi_1(0)\| P_t^{(1)}-P_t^{(2)}\|_{t^*}\left\{B_1+B_2 d_1 a+B_1^{*}\left(B_2 a+d_1 a^2/2\right)e^{\delta 1^a}\right\}$$

$$\leq \beta_1^{*}\int_o^a \left| B_1^{(1)}(\sigma)-B_1^{(2)}(\sigma)\right| do+\left(r^{*}+\|\Phi_t\|_{t*}\right)\| P_t^{(1)}-P_t^{(2)}\|_{t^*} g_1^{(a)}$$

Where $g_1(a)\rightarrow \beta_1$ as $a\rightarrow 0$. By Gronwall's inequality it will follow that

$$\left|B_1^{(1)}(a)-B_1^{(2)}(a)\right|\leq \left(r^{*}+\|\Phi_t\|_{t*}\right)\| P_t^{(1)}-P_t^{(2)}\|_{t^*}\{g_1(a)+0(a)\}$$

Using this estimate for $R_1(a)$ we get

$$|R_2(a)|\leq \left| \int_0^a [B_1^{(2)}(\sigma) B_1^{(2)}(\sigma),] b_1(a))-b_1(a-\sigma, P_t^{(1)}(a)) M_1(\sigma, a, P_t^{(1)}(a))do\right|$$

$$\left|B_1^{*}\left(r^{*}+\|\Phi_t\|_{t*}\right)\| P_t^{(1)}-P_t^{(2)}\|_{t^*}0(a) as\ a\rightarrow 0\right.$$

*Thus*

$$\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right| \leq (r^* + \| \Phi_t \|_{t*}) \| P_t^{(1)} - P_t^{(2)} \|_t o(1) as \quad a \to o$$

And hence

$$\left| \pi_1(P_t^{(1)} - \pi_1\left(P_t^{(2)}\right) \right| \leq ((r^* + \| \Phi_t \|_{t*}) \| P_t^{(1)} - P_t^{(2)} \|_t o(t) as \quad t \to o$$

Which together with a similar bound for $\left| \pi_2(P_t^{(1)} - \pi_2\left(P_t^{(2)}\right) \right|$ leads to

$$\| \pi\left(P_t^{(1)} - \pi\left(P_t^{(2)}\right) \| \leq (r^* + \| \Phi_t \|_{t*}) \| P_t^{(1)} - P_t^{(2)} \|_{t*} o(t^*)$$

As $t^* \to o$ showing that if $t^*$ is small enough, the mapping $\pi$ on

$s_{t^*}(\Phi, r^*)$ is a contraction. Hence the system (2.6)-(2,8) has a unique solution for $o \leq t \leq t^* for\ t^*$ sufficiently small and positive

Now if [0,T] is any finite interval on $R^+$, we can consider [0,T] as a finite union of intervals of length less than or equal to $t^*$ Since the estimate (2.11)is valid for any finite interval we can extend the solution of (2,6)-(2.8) form $[0, t^*]$ to [0,T] and this completes the proof,

3. Stationary age distributions and their linear stability we will now establish the existence of time independent solution

$P_t^*, P_t^*$ Of (1.1)-(1.6): such solutions satisfy

$$(3.1)\ \frac{dp\, P_2^*(a)}{da} + f_i(a, P_1^*, P_2^*)\, P_i^*(a) = 0, \quad i = 1, 2, a > 0,$$

$$(3.2)\ P_i^* = \int_o^\infty P_i^*(a) da, i = 1,\ 2,$$

$$(3,3)\ P_i^* = \int_o^\infty b_i(a,\ P_1^*,\ P_2^*) P_i^*(a) da,\ i = 1.\ \ 2$$

All solutions of (3.1)are of the from

$$(3.4)\ P_i^*(a) = P_i^*(o) \exp\left[ - \int_o^\infty f_i(s, P_1^*, P_2^*) ds \right], i = 1, \quad 2, a \geq 0,$$

Which together with (3.3) lead to

$$1 = \int_o^\infty b_1(a, P_1^*, P_2^*) \exp\left[ - \int_o^a f_1(s, P_1^*, P_2^*) ds \right] da,$$

$$(3.5)\ 1 = \int_o^\infty b_2(a, P_1^*, P_2^*) \exp\left[ - \int_o^a f_2(s, P_1^*, P_2^*) ds \right] da$$

Thus the existence of stationary solutions of (1.1)-(1.6) reduces to the existence of a pair of positive constants $P_1^*, P_2^*$ satisfying (3.5). We can now prove the following,

THEOREM 3. Assume that $f_1$ and $b_i$ satisfy ( $A_1$ ) and ( $A_2$ )and let

(3.6) $F_{i.}\ (P_1, P_2) = \int_o^\infty b_i(a, P_1, P_2) \exp\left[-\int_o^\infty f_i\,(s, P_1, P_2)ds\right] da,$

$i$ =1, 2, $P_1 \geq 0$, $P_2 \geq 0$,

**And suppose the following hold:**

(3.7) $A_{4.}$ $F_{i.}\ (0,\ 0) > 1,\ i = 1, 2$:

$A_{5.}$ *there exists a positive constant c for which*

$F_1$(C,0)< $F_{2.}$(C,0),

(3.8) $F_1$(0.C)< $F_{2.}$(0,C),

Then there exists a unique pair $(P_1^*, P_2^*)$ of real numbers $P_1^* > o$,

$P_2^* > 0$ such that

(3,9) $F_1(P_1^*, P_2^*) = 1 = F_2(P_1^*, P_2^*)$

And the unique nonegative solution of (3.1)-(3.3) is given by

(3.10) $P_2^*(a) = [P_2^* \exp[-f_i\,(s, P_1^*, P_2^*)ds])/(\int_o^\infty \exp[-\int_o^a f_i\,(s, P_1^*, P_2^*)ds]da),$

$i$ =1, 2, $a \geq 0$,

Proof Consider the elements of the surfaces defined by

Z= $F_1(P_1, P_2), z = F_2(P_1, P_2),\ P_1 \geq 0, P_2 \geq 0,$

In the ($P_1$,$P_2$,z) space. By hypothesis ($A_1$) we have $\partial F_i/\partial P_j < 0$ *and* , by ($A_2$), $\partial b_i/\partial P_j \leq 0$ for all

$P_1 \geq 0, P_2 \geq 0$, $i$ =1, 2,

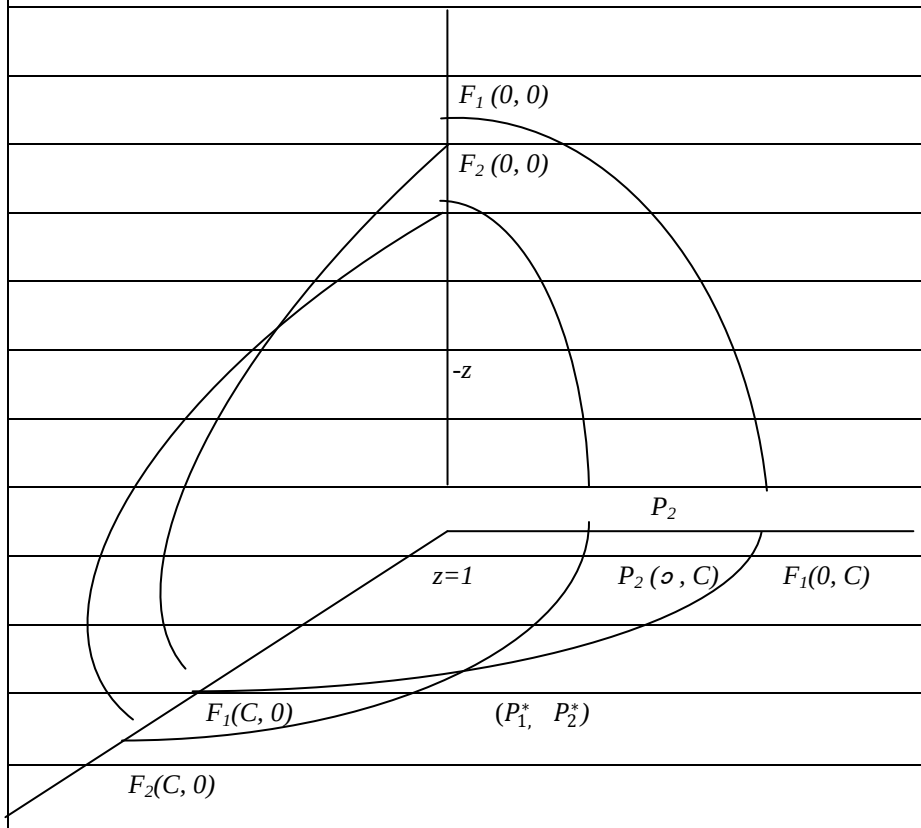
$j$ =1, 2, Hence by ($A_4$) the intersections of these surfaces with the plane z=1 define two curves on the z=1 plane connecting the I ines {z=1, $P_1 \geq 0, P_2 \geq 0,$} and {z= $1, P_2 \geq 0, P_1 \geq 0$, These two curves are defined by

$F_1(P_1, P_2 = 1$ *and* $F_2(P_1, P_2) = 1,\ F_1(P_1, P_2 \geq 0,$

Which by ($A_s$) will intersect at a unique point say $((P_1^*, P_2^*)$ with $(P_1^* > 0, P_2^* > 0$ (see figure). The uniquess of the point $((P_1^*, P_2^*)$ *is a* consequence of the smoothness of the surfaces $z = F_1(P_1, P_2)$ and $z = F_2(P_1, P_2)$

Following is a hypothetical case for Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates attributable or ascribable in the transactions and reactionary potentialities.




(The figure corresponds to the case $F_1(0,0) > F_2(0,0) > 0$. The other possibility $1 < F_1(0,0)$ is treated similarly) It will now

Follow that for such a pair $(p^*_1, p^*_2)$ the unique nonnegative solution of (3.1) - (3.3) is given by (3.10).

Now to examine the linear stability of the stationary age distributions $p^*_1, p^*_2$, we let

(3.11)
$$\begin{cases} p_i(a,t) = p_i^*(a) + u_i(a,t) \\ p_i(t) = p_i^* + p_i(t), \quad i = 1, 2, \quad a > 0, \quad t > 0, \\ p_i(t) = \int_0^\infty u_i(a,t)da, \end{cases}$$

In (1.1)-(1.5) and derive the following variational system after neglecting the nonlinear terms in the perturbations $\boldsymbol{u}_i$ and $\boldsymbol{p}_i$ $(i = 1, 2,)$ (such a procedure of linear stability analysis can be justified As has been done in Gurtin and MacCamy [4].)

(3.12)
$$\begin{cases} Du_i(a,t) = -f_i\left(a,p_1^*,p_2^*\right)u_i(a,t) - p_i^*(a)\left\{\sum_{i=1}^{2}\frac{\partial f_i}{\partial p_{\square}^*}\,p_{i\square}(t)\right\}, \\ u_i(0,t) = \int_0^\infty b_i\left(a,p_1^*,p_2^*\right)u_i(a,t)da + \int_0^\infty p_i^*(a)\left\{\sum_{\square=1}^{2}\frac{\partial b_i}{\partial p_{\square}^*}\,p_{i\square}(t)\right\}da \end{cases}$$

Where $i = 1,2$ and (3.13)
$$p_{11}(t) = p_1(t),\quad p_{12}(t) = p_2(t-\tau),$$
$$p_{21}(t) = p_1(t-\tau),\quad p_{22}(t) = p_2(t)$$

To consider the asymptotic bahaviour as $t \to \infty$ as of solutions of (3.12) we let

(3.14)
$$\begin{cases} u_i(a,t) = u\square(a)\exp[\lambda t], \quad \square(a) \to 0 \text{ as } a \to \infty, \\ p_i(t) = p_i^*\exp[\lambda t], \\ \qquad\qquad i = 1,\ 2,\ a > 0,\ t > 0, \\ p_i^* = \int_0^\infty \square(a)da, \end{cases}$$

Using (3.14) in (3.12) - (3.13),

$$\frac{d\square}{da} + \lambda\square = -f_i\left(a,p_1^*,p_2^*\right)\square(a) - p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\,p_{i\square}^*\right\},$$

(3.15)

$$\square(0) = \int_0^\infty b_i\left(a,p_1^*,p_2^*\right)\square(a)da + \int_0^\infty p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial b_i}{\partial p_j^*}\,p_{i\square}^*\right\}da,$$

Where $i = 1,\ 2,$ and

$$p_{11}^* = p_1^*,\quad p_{12}^* = p_2^*\exp\{-\lambda\tau\},$$

$$p_{21}^* = p_1^*,\ \exp\{-\lambda\tau\},\ p_{22}^* = p_2^*.$$

If we integrate the first of (3.15) and use the fact that $\square(a) \to 0$ as $a \to \infty$ we can derive that $p_1^*$ and $p_2^*$ are governed by

(3.16) $$\lambda p_1^* - \int_0^\infty p_1^*(a)\left[\left(\frac{\partial b_1}{\partial p_j^*} - \frac{\partial f_1}{\partial p_1^*}\right)p_1^* + \left(\frac{\partial b_1}{\partial p_2^*} - \frac{\partial f_1}{\partial p_2^*}\right)p_2^* e^{-\lambda\tau}\right]da$$

$$= \int_0^\infty \left[b_1(a, p_1^*, p_2^*) - f_1(a, p_1^*, p_2^*)\right]\Phi_i(a)da,$$

(3.17) $$\lambda p_2^* - \int_0^\infty p_2(a)\left[\left(\frac{\partial b_2}{\partial p_1^*} - \frac{\partial f_2}{\partial p_1^*}\right)p_1^* e^{-\lambda\tau} + \left(\frac{\partial b_2}{\partial p_2^*} - \frac{\partial f_2}{\partial p_2^*}\right)p_2^*\right]da$$

$$= \int_0^\infty \left[b_2(a, p_1^*, p_2^*) - f_2(a, p_1^*, p_2^*)\right]\Phi_2(a)da,$$

From the first of (3.15) we get

(3.18) $$\Phi_i(a) = \{\Phi_i(0)p_i^*(a)/p_i^*(0)\exp[\lambda a]\}$$
$$- p_i^*(a)\exp[-\lambda a]\int_0^a \left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\, p_i^*\Phi_i\right\}e^{\lambda s}ds, \quad i = 1, 2,$$

Which by the second of (3.15) becomes,

(3.19) $$\Phi_i(a)\left[p_i^*(a)/p_i^*(0)c_i\exp(\lambda a)\right]\left[\int_0^a p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\, p_i^*\Phi_i\right\}da\right.$$

$$\left. - \int_0^\infty b_i\left(a, p_1^*, p_2^*\right)p_i^*(a)\, e^{-\lambda a}\left\{\int_0^a\left(\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\, p_i^*\Phi_i\right)e^{\lambda s}ds\right\}da\right]$$

$$- p_i^*(a)\, e^{-\lambda a}\int_0^a\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\, p_i^*\Phi_i\right\}e^{\lambda s}ds, i = 1, 2,$$

where

(3.20) $$c_i = 11 - \frac{1}{p_i^*(0)}\int_0^\infty b_i\left(a, p_1^*, p_2^*\right)p_i^*(a)e^{-\lambda a}da, \quad i = 1, 2,$$

Now supplying $\Phi_1(a)$ and $\Phi_2(a)$ from (3.19) in (3.16)-(3.17) we derive that

$$(\lambda - A_{11})p_1^* - A_{12}p_2^*\exp(-\lambda\tau) = g_{11}(\lambda)p_1^* + g_{12}(\lambda)p_2^*\exp(-\lambda\tau),$$

(3.21)

$$-A_{21}p_1^*\exp(-\lambda\tau) + (\lambda - A_{22})p_2^* = g_{21}(\lambda)p_1^*\exp(-\lambda\tau) + g_{22}(\lambda)p_2^*$$

Where

(3.22) $$A_{ij} = \int_0^\infty p_i^*(a)\left[\frac{\partial b_i}{\partial p_j^*} - \frac{\partial f_i}{\partial p_j^*}\right]da,$$

(3.23) $g_i(\lambda) = [p_i^*(0)c_i]^{-1}\left\{\int_0^\infty \sigma_i(a)p_i^*(a)\,e^{-\lambda a}\,da\right\}$
$\left\{\int_0^\infty \left[p_i^*(a)\frac{\partial b_i}{\partial p_i^*} - b_i(a,p_1^*,p_2^*)p_i^*(a)\,e^{-\lambda a}\left(\int_0^a \frac{\partial f_i}{\partial p_i^*}\,e^{\lambda s}\,ds\right)\right]da\right\}$

$$-\int_0^\infty \sigma_i(a)p_i^*(a)\,e^{-\lambda a}\left(\int_0^a \frac{\partial f_i}{\partial p_i^*}\,e^{\lambda s}\,ds\right)da,\quad i=1,\ 2,$$

(3.24) $\sigma_i(a) = b_i(a,\ p_1^*,p_2^*) - f_i(a,\ p_1^*,p_2^*),\quad i,\ =1,\ 2.$

To solve the linear perturbational system (3.12) it is enough to find the constants $\lambda, p_1^*, p_2^*$ satisfying (3.21) since one can then use (3.19) to find the perturbations from (3.14). A nontrivial solution $(p_1^*, p_2^*)$ for (3.21) will exist if and only if $\lambda$ is a root of the equation $et\begin{bmatrix}\lambda-(A_{11}+g_{11}(\lambda)) & -(A_{12}+g_{12})e^{-\lambda\tau}\\ -(A_{21}+g_{21})e^{-\lambda\tau} & \lambda-(A_{22}+g_{22}(\lambda))\end{bmatrix}=0$

Or equivalently

(3.25) $\lambda^2 - \lambda(A_{11}+A_{22}) + A_{11}A_{22} - A_{12}A_{21}\exp[-2\lambda\tau] = \lambda s_1(\lambda) + s_2(\lambda,\tau)$

Where $S_1(\lambda) = g_{11}(\lambda) + g_{22}(\lambda)$ ,

$S_2(\lambda,\tau) = [A_{21}g_{12}(\lambda) + A_{12}g_{21}(\lambda) + g_{12}(\lambda)g_{21}(\lambda)]\exp[-2\lambda\tau]$

$-A_{11}g_{22}(\lambda) - A_{22}g_{11}(\lambda) - g_{11}(\lambda)g_{22}(\lambda)$

The following result establishes the linear local asymptotic Stability of the stationary age distributions $p_1^*(a)$ and $p_2^*(a)$

THEOREM 4 Assume the following:

$A_6.\ 0 < \beta_i^* < d_i^*\ (i,\ =1,\ 2.)\ (see(A_1)\ and\ (A_2))$;

$A_7.\ \{A_{11}+A_{22}+\|S_1(0)\|\}^2 < 4\{(A_{11}A_{22}-A_{12}A_{21}) - \|S_2(0)\|\}$,

$\|S_1(0)\| = |g_{11}(0)| + |g_{22}(0)|$ ,

$\|S_2(0)\| = |A_{11}g_{12}(0)| + |A_{12}g_{21}(0)| + |g_{12}(0)g_{21}(0)|$

$+ |A_{11}g_{22}(0)| + |A_{22}g_{11}(0)| + |g_{11}(0)g_{22}(0)|$

Then all the roots of (3.25) have negative real parts and hence the

perturbations $u_i(a,\ t) = \phi_i(a)\exp[\lambda t] \to 0\ as\ \ t\to\infty$ for $i=1,2\ and\ a\geq 0$

Proof . Define $F_1(\lambda,\tau)$ and $F_2(\lambda,\tau)$ as follows:

$F_1(\lambda,\tau) = \lambda^2 - \lambda(A_{11}+A_{22}) + A_{11}A_{22} - A_{12}A_{21}\exp[-2\lambda\tau];$

$F_2(\lambda,\tau) = \lambda S_1(\lambda) + S_2(\lambda,\tau),$

Using $(A_6)$ one can show that for $\lambda$ with re $(\lambda) \geq 0$, $c_1$ and $c_2$ are positive and bounded away from zero, Now since $A_{ij} > 0$ and $\tau \geq 0$ we find from the nature of the dependence of $g_{ij}$ on $\lambda$ that for all $\lambda$ $\mu \pm iw$ with $\mu \geq 0$ and $w \geq 0$ we have

(3.26) $$\left| F_1(\lambda,\tau) - P_2(\lambda,\tau) \right|$$

$$\geq \left| F_1(\lambda,0) \right| - F_2(\lambda,\tau) \left|\right.$$

$$\geq |\lambda|^2 - |\lambda|\{A_{11}+A_{22}+\| S_1(0) \|\} + \{A_{11}A_{22} - A_{12}A_{21} - \| S_2(0,\tau) \|\}$$

$$\geq \left[(|\lambda| - P)^2 + (q)\right]^2$$

Where

$p = \{A_{11}+A_{22}+\| S_1(0) \|\}/2,$

$q = \left\{4|(A_{11}A_{22} - A_{12}A_{21}) - \| S_2(0,T) \| | - |A_{11}+A_{22}+\| S_1(0) \| |^2\right\}/2,$

It will now follow, from (A) and (3.26), $F_1(\lambda,\tau) = F_2(\lambda,\tau)$ connot have roots with zero or positive real parts and this completes the proof.

In conclusion we remark that the condition (3.7) means that the net creation rates of each system is greater than utility; a condition of this type is well known in age wise classification of particles, atoms, molecules, cosmological phenomenon. (3.8) can be interpreted to mean that the intraspecific inhibition is higher than that of the interspecific interaction. The conditions of Theorem 4 are analytical and do not lend themselves for any worthwhile interpretation in terms of the model parameters. So also an interpretation can be recorded as a remark in respect of destruction rates of particles, properties, elements, or projections.

Note: The model is applicable to each and every module. As stated in the foregoing, there are possibilities of situation where an account itself be a receiver and it is also the one to which debit is raises. Such examples are legion in all the fields of finance, economics, political science and physics and chemistry, and of course mathematics. Classification based on time is the preeminently and predominantly suited premise for such model. We repeat that the model is to be read with each module We will expatiate the importance and implicational ramifications as we proceed with the work. Assortment-blend, collection-combinations, combo-gallimaufry, goulash-hash, hodgepodge-medley, mishmash-mixture, motley-pastiche patchwork-salmagundi, soup-stew, mixed bag-mélange, olio- olio podrida, assortment variety, Array choice, collection combination, combo diversity, hodgepodge jumble, kind medley, miscellany mishmash, mixed bag mixture, mélange potpourri, selection sort, brew concoction, beverage blend, broth compound, distillation drink, fermentation hash, hodgepodge infusion, instillation liquor, mélange miscellany, mishmash mixture, mélange potpourri, borscht bouillon, bowl brew, chowder concoction, consommé decoction, gumbo hodge-podge, olio porridge, potage potpourri, pottage puree, splash stock,

vichyssoise hodgepodge, mixture- mess, collection combination, goulash hash, jumble medley, miscellany mishmash, mixed bag mélange, olio patchwork, potpourri salmagundi, jumble clutter, disarrangement disarray, disorder farrago, gallimaufry garbage, goulash hash, litter medley, mess miscellany, mishmash mixture, muddle mélange, olio pastiche, patchwork potpourri, salmagundi scramble, shuffle snarl, tangle tumble, medley miscellany, assortment brew, collection combo, conglomeration farrago, hodgepodge jumble, melee mingling, mishmash mixed bag, mixture mélange, pasticcio pastiche, patchwork potpourri, salmagundi variety, mélange mixture, assortment combo, farrago gallimaufry, pasticcio pastiche, patchwork potpourri, salmagundi soup, stew miscellany, accumulation aggregation, anthology assortment, brew collectanea, combination combo, compilation conglomeration, accumulation diversity, farrago gallimaufry, melee mishmash, mix mixed bag, mixture muddle, mélange odds and ends, olio pasticcio, pastiche patchwork, potpourri salad, salmagundi smorgasbord, stew variety, mishmash jumble, medley mess, miscellany mixed bag, mixture mélange, olio patchwork, potpourri salmagundi, mixed bag array, assortment choice, collection combination, combo conglomerate, diversity hodgepodge, jumble medley, mélange miscellany, mishmash mixture, montage mélange, odds and ends, patchwork potpourri, smorgasbord sundries, mix/mixture assortment, combination, admixture adulteration, alloy amalgam, amalgamation assimilation, association batter, blend brew, combine combo, commixture composite, compound concoction, confection conglomeration, cross crossing dough fusion, goulash grab bag, hodgepodge hybrid, hybridization incorporation, infiltration interfusion, jumble mash, medley merger, mingling miscellany, mishmash mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, odds and ends assortment, bits and pieces, debris etcetera, hodgepodge jumble, knickknacks leavings, leftovers litter, medley melee, miscellaneous paraphernalia, miscellany mishmash, mixed bag motley, mélange notions, novelties oddments, olio particles, patchwork potpourri, rags refuse, remainder remains, remnants rest, rummage scraps, sundries sundry items, compilation mishmash, paste-up patchwork, potpourri reappropriation, reproduction synthesis, conglomeration cross, crossing dough, fusion goulash, interfusion jumble, mash medley, merger mingling, miscellany mishmash, mix mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, colluviums gallimaufry, hash mixture, pasticcio pastiche, conglomeration cross section, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multiformity multiplicity, mélange potpourri, range shift, variegation variousness, divergencies diverseness, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multifariousness multiformity, mélange potpourri, range shift, soup stew, variance variation, variegation variousness, oddment odds and ends, assortment bits, polymorphism assortment, change collection, combo conglomeration, cross section departure, discrepancy disparateness, divergencies diverseness, diversification diversity, fluctuation heterogeneity, heterogeneousness incongruity, intermixture many-sidedness, medley miscellaneousness, miscellany mishmash, mixed bag mixture, notions novelties, oddment oddments, olio particles, patchwork potpourri of the particles, interactions, accounts, transactions in the cosmos. The wisdom and understanding of our own divine essence is like a beaconing begorra, bureaucratic burdock, befuddling binary and digital digamma, that invites us inward toward ever deeper connection with all around us. Yet rather than look within, most in this world prefer to look outward for divine wisdom, thus sustaining a dependence upon a vast hierarchy that stretches between the individual and the Divine. In all our wanderings away from the Divine, humankind has obscured, obfuscated, obliterated, adjured, abnegated and abjured its most compelling features through a persistent belief systems in limitations arising out of the controlling mechanisms of this hierarchy of sub ordinates, superordinates and coordinates. Has my consciousness changed? Why that doubt? Twist and shake, ponderous novelty themes, pathos, poignancy in writing in different styles, accord-coetaneousness, coevality-coincidence, concurrence-conformity, conjunction-contemporaneousness, harmony, order, peace, simultaneousness- synchronicity exciting optimism of youth, power obmbudsman,pulsations and throbbing of pop, hard luck stories of people, modernity and mischief, grimalkins of

rurality, wobbling boats of lovers, saxophone in mouths all point to newer value orientations. The Divine dances and the “Grand design “is perpetrated outside of the confines of any hierarchical structures, be it feudal, organizational. It is complete within itself, and has a singular purpose of demonstrating the collective potential of all life within the universe. It is the archetype of perfection. It is the standard bearer of each soul's innate design and ultimate destiny. The essence of the Divine is far beyond mental conception, yet humanity's tendency is to resort to the limiting language of the hierarchical paradigm acceding- accession, accommodation- accord, accordance, acknowledging, adjustment-affiliation, affinity-alliance, amity-approving, arbitration-arrangement, assenting, authorizing, bargaining-compatibility, compliance--compromise, concert-concession, conformity-congruity, consistency-correspondence, endorsing, granting-harmony in conformity to pen chance ,predilections of the protagonist to define and understand it. Searching for Connection and Wholeness calls for not lumping elements like me made of clay feet and weak knees, owing allegiance to the fringe, without courante and courage of conviction and will for vindication, visual restlessness and exaggerated narratives of self to its own, exciting potentialities available, terrestrial passion, and skew of scams all bear an infallible demonstration for new order of consciousness. Theorists have posited a virtual smorgasbord of objects to account for dark matter, although many of them have fallen prey to observational constraints. As leading possible candidates for baryonic dark matter, there are black holes (large and small), brown dwarfs (stars too cold and faint to radiate), sun-size MACHOs, cold gas, dark galaxies and dark clusters, to name only a few. The range of particles that could constitute nonbaryonic dark matter is limited only slightly by theorists' imaginations. The particles include photinos, neutrinos, gravitinos, axions and magnetic monopoles, among many others. Of these, researchers have detected only neutrinos--and whether neutrinos have any mass remains unknown. Experiments are under way to detect other exotic particles. If they exist, and if one has a mass in the correct range, then that particle might pervade the universe and constitute dark matter. But these are very large "ifs." To a great extent, the details of the evolution of galaxies and clusters depend on properties of dark matter. Without knowing those properties, it is difficult to explain how galaxies evolved into the structures observed today. As knowledge of the early universe deepens, I remain optimistic that we will soon know much more about both galaxy formation and dark matter. What we fail to see with our eyes, or detectors, we can occasionally see with our minds, aided by computer graphics. Computers now play a key role in the search for dark matter. Historically, astronomers have focused on observations; now the field has evolved into an experimental science. Today's astronomical experimenters sit neither at lab benches nor at telescopes but at computer terminals. They scrutinize cosmic simulations in which tens of thousands of points, representing stars, gas and dark matter; interact gravitationally over a galaxy's lifetime. A cosmologist can tweak a simulation by adjusting the parameters of dark matter and then watch what happens as virtual galaxies evolve in isolation or in a more realistic, crowded universe. Old consciousness will not do. Newer one has to be imbibed and cultivated. When people are unaware of their inner wisdom and wholeness, they tend to search for order and security outside themselves. That is my personal experience. There is always search for sanctity of emotion, aesthetic expression, pursuit of harmony, unraveling realms, syncopated ending and stillness of mind and shiver of heart at the mirror broken somewhere are the strains of newer wisdom if it can be called one. Same questions shall be asked and same shall be answer in the similar vein. Uncertain of their place within the hierarchical order of the world, they define themselves based upon their insecurities. Individuals thus become only pieces of their wholeness and like shards of glass from a beautiful vase; they bear little resemblance to their aggregate beauty. Within the hierarchy, many in high positions of power have taken advantage of our collective insecurities in attempting to guide the development of all humankind. They have obscured the direct connection between the individual and the Divine through a variety of means designed to intercede between our inner essence and our divine source. Each individual can come to know them to be free of all forms of institutional and organisational and societal hierarchical control. This is not to imply that we should not trust others or join together in bonds of friendship and community. It is simply a reminder that relative truth is constantly shifting in the hands of those who desire to control. And even when the motive for their controlling behaviors may be of good will, it is still a form of control. When the revealers of "truth" within the hierarchy withhold and suppress information, they are usually positioning themselves to maintain or acquire more power

rather than to disseminate empowerment to all. The desire for connection and wholeness is a fuel that drives us to seek out and explore the hierarchical paradigm. This inner longing provides us with the motivation to seek help and guidance from a specific group within the hierarchy, and in so doing, cultivate a sense of belonging and connection. Furthermore, the hierarchical paradigm is a stage whereby we develop a sense of connection to some grand, encompassing vision. This is why the hierarchy nurtures prophets who point toward a greater vision. Spiritual leaders are able to peer deeply beneath the surface reality of life and experience how intricately connected every life form is, and how the composite of all life is intelligent far beyond measure. These visionary leaders can thus interpret reality through their personal abilities to perceive and express life's dimensional depth and limitless intelligence. Yet no one is able to articulate life's full dimensional depth and breadth with the tools of language. They can only, at best, describe their interpretation or their impressions. In actuality, all of us are able to peer beneath the surface reality of life and perceive a unique vision of the universe. We require only time and intention to develop our own interpretations. And this is precisely what many great spiritual leaders have taught. Life's deeper essence is not an absolute to be experienced by the chosen few, but an evolving, dynamic intelligence that wears as many faces as there are life forms. No group or people have the exclusive portal into the universe by which the Divine expresses itself in all its majesty. The portal is open and available to all, because the Divine is within all things. Those recognized as great prophets each produced a vision of the universe beyond what was currently defined by the hierarchy. Because their interpretations were articulated with authority and depth of insight, they became a target of debate among various groups in the hierarchy. This debate then created a polarity of belief. A sympathetic constituency emerged to defend and embellish their leader's interpretation, while established groups held it in contempt of previously held beliefs. Invariably, the leader's vision became confined and shaped into dogma by followers who desired to create a new religion or sect. Thus, this infusion of fresh insight quietly receded into the hands of the hierarchy, where its deeper meaning was obscured by the very fact that it was incorporated into a rigid structure that both protected and promoted it. A new paradigm is emerging that promotes a clear connection of individual consciousness to the compelling features of the Divine without the intervention of a hierarchy. This is when the fables and myths of history step into the light and become known as they were originally intended. This is the time when language will be transformed into a new form of communication that breaks down all barriers of control. Personal transformation, through the awakening of inner wisdom and understanding, is the pathway into wholeness. The transformation paradigm is initiated simply by the recognition that rather than the dependency-inducing ways of the hierarchical paradigm, there are accelerated, independent pathways that bypass the hierarchy and lead to self-mastery. These new pathways lead to the divine wisdom and understanding that is present within all of us. This wisdom can be accessed through the practice of three principles of transformation: seeing the Divine in all, nurturance of life, and gratitude. The application of these life principles disengages individuals from the controlling elements of the hierarchy, thus initiating the transformation experience. As there are relative truths, there are relative freedoms. As individuals evolve through the hierarchical process, an ever-increasing sense of freedom is gained, yet external forces continue to exert control through limiting language, confining belief systems, and more. These controlling influences lead to anatomical architecture, body, build, cagey carcass, construction, enclosure, fabric, flounce, form, framework, fringe, groundwork, hem, mount, mounting outline, physique, scaffold, scaffolding scheme, setting shell, stage, structure, support, system, trim, trimming, truss- valance continued reliance upon the hierarchy as it unceasingly attempts to impose a sense of inequality between us. The underlying equation of the hierarchical paradigm is: individual + hierarchy = God connection. In the case of personal transformation it is: individual + inner wisdom and understanding divine equality with all. The time has come to integrate the dominant paradigm of the hierarchy with the liberating transformation paradigm. That shall redeem all promises made to oneself. There shall be a Matrix force or balancing force that shall remain UN satisfies until conservativeness is reached. This integration occurs naturally once we have fully explored the two paradigms and develop a synthesis paradigm whereby transformation is attained by searching for connection and wholeness through our own inner wisdom and understanding. It is this combination of release from dependency on the hierarchy and transformation from within that initiates the synthesis paradigm. Once we take responsibility for our own transformation and

integration, it does not mean that the hierarchy is to be shunned or avoided. The hierarchy is quite benign as a manipulative force and merely represents one important stage in the journey toward wholeness. What is being set in motion now is the initial preparation for these paradigm shifts. More specifically, these paradigms will be simultaneously played out over the coming years. As always, it will be the choice of each one of us as to which paradigm we embrace in our journey. All of the highest imaginings of humanity are yet unaware of our deepest foundation. We have sought the upper reaches of the building, yet remain unaware of the foundation's design. It is here, at the very core of existence that the Divine is bursting forth with its creative energy while simultaneously reintegrating with its invitation to wholeness. It is here that equality is realized, not in the lofty places of relative truth lodged in the hierarchy, but rather in the deepest part of the foundational plan of life's original source and ultimate destiny. The origin and destiny of existence is the tone of equality in life. Listen for this tone – this frequency of vibration – and follow it back into the very foundation from whence all things arise and return. Consider these words as symbols only. Feel the truth that stands behind these words, and tap into this empowering energy force that reaches out for you. Know it as a tone or vibration – a resonance that waits for you around every corner in which your life will turn. It is a beacon of the Divine gathering itself into the form of language in order to usher you to a place from which you can experience the formless tone of equality – the bypass of limitation. It is the primal language of our divine source that bestows to you the freedom to generate your own deepest beauty in the expression of your highest truth.(Reference: Wikipedia)

## Gopalasamy Model For Procrastinated Conducive, Felicitous Counter Sallies In Two Specified Systems

Abstract

Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of decay rate of perturbations. Sole intention, main objective and primary concern of the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are

counted in the same vein as that of predator prey. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 3. Preliminary remarks

Almost all the equations of population ecology have been derived starting from the following simple format (Note this can be used for the development of the creed and type of systemic types in case of cosmology and astrophysics and fiancé and econometrics. To give an example of finance, we can think of a account of the company giving rise to various accounts of its employees, and that is the contribution of the single account for the development of the total system types of the given configuration). It may be account, particle, atom, molecule, proposition, problem or any such thing. We shall delineate and disseminate information and particulars as and when the model shall be applied to the problem in question.

$$\frac{dN(t)}{dt} = \{\text{an individual's contribution to population growth}\}\, N(t) \qquad (1.1)$$

where $N(t)$ denotes the size of a population of a single species (or biomass) with overlapping generations at time $t$, and $dN(t)/dt$ is the rate of change of the population size at time $t$. Subsequently one makes an assumption regarding the factor in the brackets in (1.1); in particular if one assumes that an individual entity's contribution to the population growth is denoted by a function say, f(t, N), defined suitably for all t>0, N0 then one obtains from (1.1) the so called Kolmogorov formulation in the form

$$\frac{dN(t)}{dt} = f(t, N(t))\, N(t). \qquad (1.2)$$

Various choices of $f$ together with some ecologically plausible assumptions such as the temporal constancy of the environment and density dependent effects in $f$ lead to various equations. For instance if $f(t, N) \equiv r$ (a positive constant) one obtains the Malthusian formulation in the form

$$\frac{dN(t)}{dt} = rN(t)$$

And if one assumes $f(t, N) \equiv r - (r/K)N$ for some positive constants r and K one gets the familiar logistic equation of population biology

$$\frac{dN(t)}{dt} = rN(t)\left\{1 - \frac{N(t)}{K}\right\}$$

In order to device models with time delayed effects, such effects are introduced in the average growth rate $F(t, N)$ of (1.2); for instance if $f(t, N) \equiv r - (r/K)N(t - \tau)$ for some constant $\tau > 0$ we obtain Hutchinson's[19]delay loistic equation

$$\left\{\frac{dN(t)}{dt} = rN(t)\right\}\left\{1 - \frac{N(t-\tau)}{K}\right\}. \quad (1.3)$$

If for instance

$$F(t, N) = r - (r/K)N(t) - \int_{-\infty}^{t} k(t - s)F(N(s))\,ds$$

Where $k$ is a piecewise continuous nonnegative valued function such that $\int_0^\infty k(s)\,ds = 1$ with some appropriate function $F$, we will get Volterra's model of a population which pollutes its environment and pollution has an accumulative and toxic effect; such an equation is of the form

$$\left\{\frac{dN(t)}{dt} = N(t) = \right\}\left\{r - (r/K)N(t) - \int_{-\infty}^{t} k(t - s)F(N(s))\,ds\right\} \quad (1.4)$$

Which has been studied in different formats by several authors (Cushing [12], Cohen *et al.* [10], and Gopalasamy [16])

One of the significant advantages of deriving various models from the prototype (1.2) is that if the initial; conditions (initial population sizes) are nonnegative then the nonnegativity of the population size $N(t)$ for $t > 0$ follows from the fact that solutions of (1.2) satisfy

$$N(t) = N(0) \exp\left[\int_0^1 f\left(s, N(s)\right)ds\right]$$

Which implies that $N(t) \neq 0$ if $N(0) \neq 0$.

While the time delayed logistic equation (1.3) explains some of the observed oscillations of certain populations, its derivation has not been entirely satisfactory. For instance it has been pointed out by Ricklefs ([27], page 488) that "time lags might be expected to occur primarily in stabilizing processes involving production rather than death because death is an immediate response to environmental change". Recently Cushing [13] has examined the effects of response and maturation delays in a number of single and multi system models; as it has been noted by Cushing [13], the models involving delayed reproduction or not necessarily derivable from the Kolmogorov type systems of the form (1.2) by suitably defining the average growth rate $(1/N)\,(dN/dt) = f(t, N(t))$. While Cushing [13] has been mostly concerned with the possibility of delay induced oscillations we will consider a class of delays in reproduction and interspecific interaction which maintain the stability of the systems.

Our starting point to drive models with delayed reproduction is the following balance equation which assumes that there is no immigration or emigration. Account, particle, atom, molecule, proposition, problem et al

$$\frac{dN(t)}{dt} = \text{birth rate} - \text{death rate.}$$

(1.5)

Now for instance if we consider a population of adult flies then the birth rate at time $t$ depends on the adult population at time $t - \tau$ where $\tau$ is the time required for the larvae to become adult flies. If the birth and death rates are governed by density dependent factors then we have from (1.5) the following time delayed model

$$\frac{dN(t)}{dt} = b\left(N(t-\tau)\right) - m(N(t))$$

(1.6)

Where the functions $b$ and $m$ denote the density dependent birth and death rates of Account, particle, atom, molecule, proposition, problem et al respectively. If the time delay is continuously distributed then one can consider in the place of (1.6) an equation of the form

$$\frac{dN(t)}{dt} = b\int_{-\infty}^{t} k\left(t-s\right)N(s)ds - m\left(N(t)\right)$$

Or equivalently

$$\left\{ \frac{dN(t)}{dt} = b \int_0^{\infty} k(s) N(t-s) ds - m(N(t)) \right. \quad (1.7)$$

Where the delay kernel *k* denotes the weight of the delayed of the delayed effects.

It is not always obvious that suitable nonnegative initial conditions for (1.6) and (1.7) will imply the nonnegativity of the solutions of (1.6) and (1.7). The purpose of this article is to show that in two species systems of competitive and prey- predator communities, delayed reproduction and delayed interactions donot destabilise the systems if the self- regulating negative feedback effects are sufficiently strong compared to interspecific interactive effects and if the self regulating effects are realised with no time delays. When we say that time delays do not destabilise, we will mean that delays do not render an asymptotically (locally) stable steady state unstable in such a manner that the loss of stability either leads

To a delay – induced bifurcation to persistent and undamped oscillations (as in the case of Hopf- bifurcation) or the relevant linear variational system has unbounded solutions; that is a mathematical instability is not induced by delays in the sense, that the roots of the characteristic equation are pushed to the imaginary axis or to the right of it on the complex plane. However delays can make an otherwise stable system less stable which means that the rate of convergence to zero of solutions of the variational system can decrease with an increase of delay and the variational system can even oscillate (perhaps violently) before converging to the trivial state. Such a reduction in stability can happen even when all the roots of the characteristic equation have negative real parts. For our purposes such a reduced stability is still stability. In order to see the effects of delays on such less stable systems we will propose below a method of estimating the rate of decay of local perturbations based on the variational system and indirectly relying on the characteristic equation. For a special class of two species models with a single delay we have obtained (Gopalsamy [16]) estimates of decay rates (or return times) usingthe characteristic equation; such a procedure is almost impossible in the case of many species systems with several delays. The method we propose in Section 6 below for the estimate of decay rates can be adapted to *n*- species systems with $n^2$ discrete delays.

Let $x_1(t)$ and $x_2(t)$ denote the population sizes (or biomasses) of two systems competing for a common pool of resources with common pool of resources.. Let $b_i$ and $m_i$ $(i = 1, 2)$ denote the respective density dependent birth (potential) and death rates so that in the absence of any time delays the population sizes are governed by (of

Account, particle, atom, molecule, proposition, problem et al). In addition as usual, we assume that the system studied follow Haimovici's equations, for which the dove tailing explanation has been appended

$$\frac{dx_1(t)}{dt} = b_1\left(x_1(t)\right) - m_1(x_1(t)), x_2(t))$$

$$\frac{dx_2(t)}{dt} = b_2\left(x_2(t)\right) - m_2(x_1(t)), x_2(t))$$

(2.1)

In order to make the system (2.1) represent a competition system we will make the following assumptions on the birth and death rates; $_1$, $m_1$ ($i$ = 1, 2) are continuous with continuous partial derivatives for all $x_1$ 0 ($I$ = 1, 2); also we assume (see for instance Rescigno and Richardson [26]) $\frac{\partial b_i}{\partial x_1} > 0, \frac{\partial m_1}{\partial x_j} > 0$ for $x_i > 0$ ($i, j$ = 1, 2), (2.2)

(iv) $b_1(0) = 0$, $m_1(0, x_2) \equiv 0, b_2(0) = 0$, $m_2(x_1, 0) \equiv 0$,

(2.3)For some $x^*_{1,} > 0, x^*_{2,} > 0$ we have

$$b_1(x^*_1, -m_1(x^*_1, 0) = 0,$$

$$b_2(x^*_2, -m_1(0, x^*_2) = 0,$$

(2.4)

(v) There exist positive constants $\delta_1$, $\delta_2$ such that

$b_1(\delta_1) - m_1(\delta_1, x_2) < 0$ ,

$b_2(\delta_2) - m_2(x_1, \delta_2) < 0$, for $x_1$ 0, $x_2$ 0, (2.5)

(vi) There exists a point $(\alpha, \beta), \alpha > 0, \beta > 0$ such that

$$b_1(\alpha) - m_1(\alpha, \beta) = 0 = b_2(\beta) - m_2(\alpha, \beta).$$

(2.6)

The conditions on $b_i$ in (2.2) mean that the potential (not average) birth rates are density dependent and any negative density dependent feedbacks to the opened or birth rate of Account, particle, atom, molecule, proposition, problem et al rates are included in the destruction rates; the conditions on $m_i$ indicate intraspecific and interspecific confederation needs and necessities. (2.3) implies that (0, 0) is a (trivial) steady state of the system (2.1) as it is in

the case of all ecological models; (2.4) means that in the absence of either of the two species the other one has a nontrivial steady state; (2.5) implies that each of the two populations cannot grow unboundedly since when $x_1(t) > \delta_i, \frac{dx_1}{dt} < 0$. (2.6) means that there exists a nontrivial steady state $(\alpha, \beta)$ in the interior of the positive quadrant of the $(x_1, x_2)$- plane. These assumptions are standard and are minimal in the opinion of the author.

An immediate example of (2.1) –(2.6) is the familiar Volterra –Lotka competition system described by

$$\left.\begin{aligned}\frac{dx_1(t)}{dt} &= r_1x_1(t) - a_{11}x_1^2(t) - a_{12}x_1(t)x_2(t),\\ \frac{dx_2(t)}{dt} &= r_2x_2(t) - a_{21}x_1(t)x_2(t) - a_{22}x_2^2(t),\end{aligned}\right\} \quad (2.7)$$

Where $r, a_{ij}, (i, j = 1,2)$ are positive constants satisfying the condition

$$\frac{a_{11}}{a_{21}} > \frac{r_1}{r_2} > \frac{a_{12}}{a_{22}}.$$

One of the first questions to be resolved for (2.1) is the following: if $x_i(0) > 0$, $x_2(0) > 0$ does it follow that the corresponding solutions of (2.1) satisfy the condition $x_1(t) > 0$, $x_2(t) > 0$ for those $t$ for which a solution of (2.1) exists? In the $(x_1, x_2)$ phase plane of (2.1) we have $(0, 0), (x_1^*, 0), (0, x_2^*), (\alpha, \beta)$ as the steady states of (2.1) and also it will follow from (2.3) that the co ordinate axes

$x_1 = 0$ and $x_2 = 0$ are invariant sets f (2.1) ; such invariance of the coordinate axes will imply that if $x_1(0) > 0$, $x_2(0) > 0$ then the corresponding population trajectory $(x_1(t), x_2(t))$ of (2.1) cannot reach the outside of the nonnegative quadrant $x_1 \geq 0$, $x_2 \geq 0$.

Our plan is to derive first a set of sufficient conditions for the coexistence of the two competing species and then in the next section show that under the same conditions the two species can coexist if there are time delays in the reproduction and interspecific interactions. We recall that the two competing species described by (2.1) are said to coexist if and only if the nontrivial steady state $(\alpha, \beta)$ is (at least locally) asymptotically stable. To examine such a stability of $(\alpha, \beta)$ we derive the corresponding linear variational system for the perturbations $X_1, X_2$ where

$$x_1(t) \equiv \alpha + X_1(t), \qquad x_2(t) \equiv \beta + X_2(t)$$

Such a variational system is found to be

$$\left.\begin{aligned}\frac{dX_1(t)}{dt} &= \beta_{11}X_1(t) + \beta_{12}X_2(t),\\ \frac{dX_2(t)}{dt} &= \beta_{21}X_1(t) + \beta_{22}X_2(t),\end{aligned}\right\} \quad (2.8)$$

Where

$$\beta_{1J} = \left(\frac{\partial b_1}{\partial x_J} - \frac{\partial m_1}{\partial x_j}\right) \text{ evaluated at } (\alpha, \beta) \qquad (i, j = 1,2).$$

(2.9)

The steady state $(\alpha, \beta)$ of (2.1) is locally asymptotically stable if and only if the trivial solution (0, 0) of (2.8) is asymptotically stable. We can derive the following result from elementary considerations.

In the system (2.1)-(2.6) assume the following:

$$\left.\begin{aligned}\frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1},\\ \frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1},\end{aligned}\right\} \text{ at } (\alpha, \beta)$$

(2.10)

Then the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (2.1) – (2.6).

The proof follows easily if we note that the characteristic equation associated with the linear system (2.8) is given by

$$\text{Det}\begin{pmatrix}\lambda-\beta_{11} & -\beta_{12}\\ -\beta_{21} & \lambda-\beta_{22}\end{pmatrix} = 0$$

Or equivalently,

$$\lambda^2 - \lambda(\beta_{11} + \beta_{22}) + \beta_{11}\beta_{22} - \beta_{12}\beta_{21} = 0. \quad (2.11)$$

It follows from (2.2) and (2.10) that,

$$\beta_{11} < 0, \beta_{11}\beta_{22} - \beta_{12}\beta_{21} > 0$$

And hence the roots of (2.11) have negative real parts implying the asymptotic stability of the trivial solution of (2.8)

If apply the condition (2.10) to the Volterra-Lotka model (2.7) we note that when a nontrivial steady state say $(\alpha, \beta)$ exists for (2.7), (2.10) leads to

$a_{11}\alpha > a_{21}\,\beta \; and \; a_{22}\,\beta > a_{12}\,\alpha,$

Which together imply $a_{11}\,a_{22} > a_{12}a_{21}$. If $(\alpha, \beta)$ exists it is known that $a_{11}\,a_{22} > a_{12}a_{21} > 0$ is a sufficient condition for the (local) asymptotic stability of $(\alpha, \beta)$.

The condition (2.10) which guarantees the coexistence of competing populations in (2.1)-(2.6) simply means that the intraspecific self – regulating feedback effects $\partial m_i / \partial x_1 \; (i = 1, 2)$ bare higher than its own positive feedback $\partial b_1 / \partial x_1 \; (i = 1, 2)$ together with its negative (competitive) effects on its competitor $\partial m_I / \partial x_J \; (i, \; j = 1, 2; i \neq j \;)$.

## 4. Coexistence under delayed production and interactions, transactions

We will now establish that two competing species with a sufficiently strong intraspecific regulation can coexist in the presence of delayed reproduction and delayed interspecific interactions. It is well known that if there are delays in the self – regulating intraspecific effects, stationary coexistence may not always be possible (Gopalsamy [16], Gopalsamy and aggarwala [15], Shibata and Saito [28]). Delayed reproduction is more general and widespread in most biological populations. To build a competition model with such delays we proceed as follows.

Let $\tau_{ij} \; (i, j = 1, 2)$ be a set of nonnegative constants with $\tau = \max \tau_{ij} \; , i, j = 1, 2,$ and suppose that the two ompeting species display delayed reproduction and delayed interspecific interactions with no delays in the self – regulating negative feedback effects. Such a competition system in a constant environment can be modeled by the autonomous delay differential system

$$\frac{dx_1(t)}{dt} = b_1\big(x_1(t - \tau_{11})\big) - m_1\big(x_1(t), x_2(t - \tau_{12})\big),$$

$$\frac{dx_2(t)}{dt} = b_2\big(x_2(t - \tau_{22})\big) - m\big(x_1(t - \tau_{21}), x_2(t)\big), \qquad (3.1)$$

Where both the creation and destruction rates and the morality rates $b_1$ and $m_1 \; (i = 1, 2)$ satisfy the same conditions (2.2)-(2.6) as in the previous section. Along with (3.1) we will

Suppose that the initial population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) sizes are provided in the form

$$x_i(s) = \phi_i(s) \; \text{0}, \qquad s \in [-\tau, 0], \qquad i = 1,2. \qquad (3.2)$$

(Note that the specification of $\phi_i$ need not be on the entire interval $[-\tau, 0]$; but such a hypothesis is convenient.) Since (3.1) is not of the standard Kolmogorov-type, we have to establish that the solutions of (3.1) will remain nonnegative for nonnegative initial conditions.

To prove the nonnegativity of solutions of (3.1)-(3.2), let $\left\{\psi_1^{(n)}(s)\right\}, s \in [-\tau, 0], i = 1,2; (n = 1,2, \ldots)$ be a sequence of strictly positive continuous functions such that

$$\lim_{n\to\infty} \psi_1^{(n)}(s), = \psi_1 \; (s), \qquad s \in [-\tau, 0], \qquad i = 1,2 \qquad (3.3)$$

Let $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ be the solution of (3.1) corresponding to the initial condition

$$\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \left\{\psi_1^{(n)}(s), \psi_2^{(n)}(s)\right\}, \quad s \in [-\tau, 0], \tag{3.4}$$

Consider the solution $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ for $t \in [0, \tau]$ (assuming it exists on the whole of $[0, \tau]$). Suppose $x_1^{(n)}$ does not remain nonnegative for all $[0, \tau^*]$; then there exists a $t^* = [0, \tau^*](\tau^* =$ the positive minimum of $\tau_{ij}$ $(i, j = 1, 2))$ for which

$$x_1^{(n)}(t^*) = 0 \ and \ x_1^{(n)}(t)(t^* - \tau_{11})) > 0.$$

It will follow from (3.1) and the positivity of the initial condition that

$$\frac{dx_1^{(n)}(t^*)}{dt} = b_2\left(x_1^{(n)}(t^* - \tau_{11})\right) > 0 \tag{3.5}$$

Where we have used the properties of $b_i$ and $m_i$ from (2.2)-(2.6); it is found that (3.5) contradicts the definition of $t^*$. Thus $x_1^{(n)}(t) > 0$ on $[0, \tau^*]$; similarly $x_2^{(n)}(t) > 0 \ on \ [0, \tau^*]$. We can repeat this procedure for intervals of the form $[\tau^*, 2\tau^*], [2\tau^*, 3\tau^*] \ etc$. Thus it will follow that so long as $(x_1^{(n)}(t), x_2^{(n)}(t))$ exists, we have

$$x_1^{(n)}(t) \geq 0 \ x_2^{(n)}(t) \geq 0.$$

If we now consider the limits as $n \to \infty$, we get

$$\lim_{n\to\infty}\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \{x_1(t), x_2(t)\} \tag{3.2}$$

With $x_1(t)$ ,, $x_2(t)$ where $(x_1(t), x_2(t))$ is the solution of (3.1)-(3.2) and this is a consequence of the continuous dependence of solutions of (3.1) on initial conditions (Hale [17], page 41).

A second question for (3.1) with (2.2)-(2.6) is whether solutions of (3.1) exist for all $t > 0$? suppose now that solutions of (3.1) do not exist for all $t$ ; then there exists a $t_1 > 0$ such that at least for one of either of $x_1 \ or \ x_2$ we have

$$\lim_{t\to t_1-} x_i(t) = \infty. \tag{3.7}$$

To be specific let us suppose that $\lim_{t\to t_1-} x_i(t) = \infty$; let $t_2$ be the first instant of time for which $x_1(t_2) = \delta_1, t_2 < t_1$. It will follow from (3.1) that

$$\frac{dx_1(t_2)}{dt} = b_1(x_1(t_2 - \tau_{11})) - m_1(x_1(t_2), x_2(t_2 - \tau_{12}))$$

$$< b_1(x_1(t_2)) - m_1(x_1(t_2), x_2(t_2 - \tau_{12}))$$

$$< b_1(\delta_1) - m_1(\delta_1, x(t_2 - \tau_{12})) < 0 \quad (3.8)$$

Which contradicts the definition of $t_2$. Thus solutions of (3.1), (2.2) – (2.6) exist for all $t \geq 0$.

We are now ready to consider the effects of time delays in reproduction and interspecific interactions on the coexistence of the two competing species. The steady state $(\alpha, \beta)$ of (2.1) is again a steady state of (3.1). To examine the (local) asymptotic stability of $(\alpha, \beta)$ in (3.1) we let

$$x_1(t) = \alpha + Y_1(t), \quad x_2(t) = \beta + Y_2(t)$$

In (3.1) and derive that the linear variational system governing the perturbations $Y_1, Y_2$ is given by

$$\frac{dY_1(t)}{dt} = -m_{11}Y_1(t) + b_{11}Y_1(t - \tau_{11}) - m_{12}Y_2(t - \tau_{12}),$$

$$\frac{dY_2(t)}{dt} = -m_{22}Y_2(t) + b_{21}Y_1(t - \tau_{21}) - m_{22}Y_2(t - \tau_{22}),$$

Where

$$m_{ij} = \frac{\partial m_1}{\partial x_j}, \quad b_{ii} = \frac{\partial b_i}{\partial x_i} \; at \; (\alpha, \beta). \quad (3.10)$$

It is well known (Bellman and Cooke [3], page 336) that the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) if the trivial solution of the variational system (3.9) is asymptotically stable. We formulate our first result in the following form whose proof is given in the appendix.

*THEOREM 3.1. Assume that in the two systemic types competition model system (3.1) with delayed density dependent birth and death rates $b_i$ and $m_i$ $(i = 1, 2)$ satisfy the conditions (2.2)-(2.6). Suppose that the conditions (2.10) hold. Then the nontrivial steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) whatever the delays be $\tau_{ij} \geq 0$ $(i = 1, 2)$.*

Let $y_1(t) and\ y_2(t)$ respectively denote the population sizes (or biomasses) of a prey species and a predator species which feeds exclusively on the single prey species in a temporally constant environment. In the format of (1.5) the following system will describe the rates of change of the two population sizes:

$$\frac{dy_1(t)}{dt} = f_1(y_1(t)) - g_1(y_1(t), y_2(t)),$$

$$\frac{dy_2(t)}{dt} = f_2(y_1(t)) - y_2(t) - g_2(y_2(t)), \quad (4.1)$$

Where $f_i\ and\ g_i\ i = 1,2)$ respectively denote the birth and death rates of the two species; in order to make (4.1) denote the dynamics of a prey – predator species we will assume that $f_i\ and\ g_i$ satisfy the following conditions (see for instance Rescigno and Richardson [26]); $f_i\ and\ g_i$ are continuous with continuous partial derivatives for all $y_i \geq 0$ $(i = 1, 2); also\ we\ require$

$$\frac{\partial f_1}{\partial y_1} > 0; \quad \frac{\partial g_1}{\partial y_1} \; 0; \quad \frac{\partial g_1}{\partial y_2} \; 0,$$

$$\frac{\partial f_2}{\partial y_1} \; 0; \quad \frac{\partial g_2}{\partial y_2} \; 0; \quad \frac{\partial g_2}{\partial y_2} \; 0, \quad \text{Where } y_1 \; 0, y_1 \; 0, \tag{4.2}$$

(v)

$$f_1(0) = 0, \quad g_1(0, y_2 \equiv 0,$$

$$f_2(0, y_2) = 0 = f_2(y_1, 0); \quad g_2(0) = 0, \tag{4.3}$$

(vi) For some constant $y_1^* > 0$ we have

$$f_1(y_1^*) - g_1(y_1^* 0) = 0, \tag{4.4}$$

(vii) There exists a point $(\alpha^* \beta^*), \alpha^* > 0, \beta^* > 0$ such that

$$f_1(\alpha^*) - g_1(\alpha^*, \beta^*) = 0,$$

$$f_2(\alpha^*, \beta^*) - g_2(\beta^*) = 0, \tag{4.5}$$

(viii) There exist positive constants $n_1, n_2$ such that

$$f_1(\eta_1) - g_1(\eta_1, y_2) < 0, y_2 \; 0,$$

$$f_2(\eta_1, \eta_2) - g_2(\eta_2) < 0 \tag{4.6}$$

The conditions (4.2) mean that the system (4.1) is of the oppugnacy systems type; (4.3) will imply that (0,0) is a trivial steady state of (4.1); (4.4) means that in the absence of the predator, the prey species has a nontrivial (positive) steady state;

(4.5) implies that (4.1) has a nontrivial steady state in the interior of the positive quadrant of the state species of (4.1); (4.6) can be shown (see below) to imply that no credit type and debit type system types can grow unboundedly.

An example of (4.1)-(4.6) is the familiar Volterra-Lotka prey-predator system

$$\frac{dy_l(t)}{dt} = r_l y_l(t) - [a_{ll} y_l^2(t) + a_{12} y_l(t) y_2(t)],$$

$$\frac{dy_2(t)}{dt} = a_{12} y_1(t) y_2(t) - [r_2 y_2(t) + a_{22} y_2^2(t)],$$

Where $r_l, a_{ij}\,(i, j = 1,2)$ are positive costants such that

$$a_{21} r_1 - r_2 a_{11} > 0.$$

As in the previous case of competition we have to verify that when $y_1(0) > 0, y_2(0) > 0$ the corresponding solution of (4.1) will be nonnegative for $t \; 0$ for which the solution exist. Since the co-ordinate axes $y_1 = 0, y_2 = 0$ are invariant sets for (401) by virtue of (4.3), it will follow that the solution of (4.1) starting from the interior of the positive quadrant cannot reach the outside of that quadrant. This is a consequence of the invariance of the co-ordinate axes for (4.1)

To prove that solutions of (4.1)-(4.6) exist for all t 0 let us suppose that $y_l 0 \; and \; y_2(0) satisfy \; 0 < y_l(0)$ satisfy $0 < y_l(0) < \eta_l, 0 < y_2(0) < \eta_2$ ( on the other hand if either one of $y_l(0) \; \eta_l and \; y_2(0) \; \eta_l$ or both hold, it will follow that the corresponding $y_l(t)$ will decrese by virtue of (4.6) at least for small $t \; 0,$ ). Suppose $y_l(t)$ does not

exist for all $t$ ; then for some $t_l < \infty\ lim_{t\to t_l-} y_l(t) = \infty$and let $t_2 > 0$ be the first instant for which $y_1(t_2) = \eta_1$; then we have from (4.1) and (4.6) that

$$\frac{dy_1(t_2)}{dt} = f_1(\eta_1) - g_1(\eta_1, y_2) < 0\ (assumimg\ y_1(0) < \eta_1)$$

Which contradicts the definition of $t_2$ showing that $y_1(t)$ exists for all $t$ ; if $y_2(t)$ is not defined for all $t$ , then for some $t_3$ we will have $lim_{t-t_3-} y_2(t) = \infty$;now let $t_4$ be the first instant for which $y_2(t_4) = \eta_2$; then from the properties of $f_2$ we have

$$\frac{dy_2(t_4)}{dt} = f_2(y_1(t_4), \eta_2) - g_2(\eta_2)$$ $$f_2(\eta_1, \eta_2) - g_2(\eta_2)$$

$$< 0 (by\ (4.6) assuming\ 0 < y_1(0) < \eta_1, 0 < y_2(0) < \eta_2)$$

Which again contradicts the definition of $t_4$. Thus $y_2(t)$ is defined for all $t$ .

We will now derive a set of sufficient condition for the (local) asymptotic stability of the steady state $(\alpha^*, \beta^*)$ of (4.1). in the next section we will show that under the same conditions and in the presence of various delays , the steady state $(\alpha^*, \beta^*)$remains (locally) asymptotically stable.

To examine the asymptotic stability of the nontrivial steady state $(\alpha^*, \beta^*)$ in (4.1) we let

$$y_1(t) = \alpha^* + U_1(t), \qquad y_2(t) = \beta^* + U_2(t)$$

In (4.1) and derive the liner variational system in the perturbations $U_1, U_2$ as follows:

$$\frac{dU_1(t)}{dt} = \frac{\partial f_1}{\partial y_1} U_1(t) - \frac{\partial g_1}{\partial y_1} U_1(t) - \frac{\partial g_1}{\partial y_2} U_2(t),$$

$$\frac{dU_2(t)}{dt} = \frac{\partial f_2}{\partial y_1} U_1(t) - \frac{\partial f_2}{\partial y_2} U_2(t) - \frac{\partial g_2}{\partial y_2} U_2(t), \qquad (4.7)$$

Where all the partial derivatives in (4.7) are evaluated at $(\alpha^*, \beta^*)$. The steady state $(\alpha^*, \beta^*)$ of (4.1) is (locally) asymptotically stable if and only if the trivial solution of (4.7) is asymptotically stable which will be the case if and only if the roots of the characteristic equation

$$\text{Det} \begin{pmatrix} \lambda - (f_{11} - g_{11}) & g_{12} \\ -f_{21} & \lambda - (f_{22} - g_{22}) \end{pmatrix} = 0,$$

Where $f_{ij} = \frac{\partial f_1}{\partial y_l}$, $g_{ij} = \frac{\partial g_1}{\partial y_j}$, at $(\alpha^*, \beta^*)$ (4.8)

( i,j =1, 2 )

Have negative real parts. If $\lambda_1, \lambda_2$ are the riits of (4.8) then

$$\lambda_1 + \lambda_2 = (f_{11} - g_{11}) + (f_{22} - g_{22}),$$

$$\lambda_1\lambda_2 = (f_{11} - g_{11})(f_{22} - g_{22}) + f_{21}g_{12}. \quad (4.9)$$

The proof of the following result is an immediate consequence of (4.7)-(4.9)

THEOREM 4.1 *In the CREDIT-DEBIT system (4.1)-(4.6) assumes that the partial derivatives satisfy the following additional conditions:*

$$\frac{\partial g_1}{\partial y_1} > \frac{\partial f_1}{\partial y_1} + \frac{\partial f_2}{\partial y_1},$$

$$\frac{\partial g_2}{\partial y_2} > \frac{\partial f_2}{\partial y_2} + \frac{\partial g_1}{\partial y_2}, \quad \text{Evaluated at } (\alpha^*, \beta^*), \quad (4.10)$$

*Then the steady state $(\alpha^*, \beta^*)$, is (locally) asymptotically stable for the prey- predator system.*

## 6. Delays in reproduction, predation and gestation

A number of authors have proposed and analysed several models of prey- predator systems with time delays. Most of the models have been derived from the so called models of the Kolmogorov-type f (1.2) with time delays incorporated in the average growth rates of the prey and predator species. Volterra [32] proposed the following model for a prey-predator system. The equations are used for the formulation of Credit-debit system of atoms, accounts, economic variable, cosmological variable, astrophysical; variables, and qubit systems for which the model is being applied.

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \delta_1 N_1(t) - \int_{\infty}^{t} F_1(t-s)N_2(s)ds],$$

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 + \delta_2 N_1(t) - \int_{\infty}^{t} F_2(t-s)N_1(s)ds], \quad (5.1)$$

Where $r_1, r_2, \delta_1, \delta_2$ *are positive constant and* $F_1, F_2$ is nonnegative continuous delay kernel defined and integrable on $[0, \infty]$. (5.1) has no predator self-regulating negative feedback. Brelot [6] considered a modified format of (5.1) in the form

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \lambda_1 N_1(t) - \delta_1 N_2(t) - \int_{-\infty}^{t} F_1(t-s)N_2(s)ds], \quad (5.2)$$

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 - \lambda_2 N_2(t) - \delta_2 N_1(t) - \int_{-\infty}^{t} F_2(t-s)N_1(s)ds],$$

Which has term representing intraspecific competition among both predator and prey species. Under suitable condition on the various parameters in (5.2) one can show that (5.2) has a (locally) asymptotically stable steady state. A number of intergrodifferential equation models of the types (5.1) and (5.2) have been investigated by Cushing [12]

Starting from Hutchinson's (1948)[19] delay logistic equation, May[23] has proposed the following system

$$\left.\begin{aligned} \frac{dN_1(t)}{dt} &= rN_1(t)\left[1 - \frac{N_1(t-\tau)}{K}\right] - \alpha N_1(t)N_2(t), \\ \frac{dN_2(t)}{dt} &= -bN_2(t) + \beta N_2(t)N_1(t)\,, \end{aligned}\right\} \quad (5.3)$$

Where r, $\tau$, K, $\alpha, \beta, b$ are positive constants; (5.3) contains a single discrete delay; one can modify (5.3) and incorporate a continuously distributed delay in it so that (5.3) can be written as (May [23])

$$\left.\begin{aligned} \frac{dN_1(t)}{dt} &= rN_1(t)[1 - \int_{-\infty}^{t} F(t-s)N_1(s)ds] - \alpha N_1(t)N_2(t), \\ \frac{dN_2(t)}{dt} &= -bN_2(t) + \beta N_1(t)N_2(t). \end{aligned}\right\} \quad (5.4)$$

The first model of a prey-predator type which departs from the Kolmogorov-type formulation is due to Wangersky and Cunningham [33] who considered the model system

$$\left.\begin{aligned} \frac{dN_1(t)}{dt} &= a_1N_1(t) - b_1N_1^2(t) - c_1N_1(t) - N_2(t), \\ \frac{dN_2(t)}{dt} &= a_2N_2(t) + c_2N_1(t-\tau)N_2(t-\tau), \end{aligned}\right\} \quad (5.5)$$

Where $a_1, a_2, b_1, c_1, c_2, \tau$ are positive constants; (5.5) means that a duration of $\tau$ time units elapses when an individual prey is killed and the moment when the corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the

case may be( Note the author deals for predator – population.) Recently Cushing [13] has considered a model of the form

$$\frac{dN_1(t)}{dt} = rN_1(t)\left[1 - \frac{N_1(t)}{k} - \alpha N_2(t)\right],$$

$$\frac{dN_2(t)}{dt} = -d_2^*(t) + bN_1(t)\int_0^\infty \delta(a)N_2(t-a)e^{-d*a}\,da, \quad (5.6)$$

Where $r, k, \alpha, d^*, b$ are positive constants and $\delta$ is related to the age dependent fecundity of the predator species.

With the models (5.1)-(5.6) in the background let us now consider a time delayed model similar to (4.1) in the following form:

$$\frac{dN_1(t)}{dt} = f_1\big(N_1(t-\tau_{11})\big) - g_1\big(N_1(t), N_2(t-\tau_{12})\big),$$

$$\frac{dN_2(t)}{dt} = f_2\big(N_1(t-\tau_{21})\big), N_2(t-\tau_{12}) - g_1(N_2(t)), \quad (5.7)$$

Where $\tau_{ij}$ $(i,j=1,2)$ are nonnegative constants with $\tau = \max \tau_{ij}, i,j = 1,2$; the birth and death rate functions $f_1, f_2, g_1, g_2$ satisfy the same conditions as in (4.1)-(4.6). Along with (5.7) we specify the initial conditions as follows:

$$N_1(s) = \phi_1(s) > 0, \quad s \in [-\tau, 0].$$

$$N_1(s) = \phi_1(s) > 0,$$

We3 will assume that in addition to (4.2) we have for (5.7) that

$$\frac{\partial f_2}{\partial N_1} > 0, \frac{\partial f_2}{\partial N_2} > 0, f_2(N_2, N_2) > 0 \; whenever \; N_1 > 0, \; N_2 > 0. \quad (5.9)$$

It is not difficult to show by the methods previously used for the competition system that whenever $N_1(s) > 0$ and $N_2(s) > 0$ for $s \in [-\tau, 0]$, solution of (5.7) exist for all $t \geq 0$ and remain nonnegative for all $t \geq 0$. also $(\alpha^*, \beta^*)$ of (4.5) is a steady state of (5.7).

We will discuss the form and meaning of the delays in (5.7) in the next section. The effects of the various delays in (5.7) on the (local) asymptotic stability of $(\alpha^*, \beta^*)$ can be derived from the following theorem whose proof is quite similar to that of Theorem 3.1 and hence the details of the proof are omitted.

THEOREM 5.1 *In the time delayed prey-predator system(* corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) *) (5.7) let the conditions (4.2)-(4.6), (5.8) and (5.9) hold for the birth and death rates; furthermore suppose the condition (4.10) hold. Then for all nonnegative delays in (5.7), its nontrivial steady state* $(\alpha^*, \beta^*)$ *is (locally) asymptotically stable.*

## 6. Estimation of decay rates

We will first derive a convention sufficient condition for the asymptotic stability of the trivial solution of a linear system of the form

$$\left.\begin{aligned}\frac{dY_1(t)}{dt} &= -m_{11}\,Y_1(t) + b_{11}Y_1(t-\tau_{11}) - m_{12}Y_2(t-\tau_{12}),\\ \frac{dY_2(t)}{dt} &= -m_{22}\,Y_1(t) + m_{21}Y_1(t-\tau_{21}) - b_{22}Y_2(t-\tau_{22}),\end{aligned}\right\} \quad t>0, \quad (6.1)$$

Where $m_{1J}$ , $b_{i1}$ $(i,j=1,2)$ are positive constants while $\tau_{1J}(i,j=1,2)$ are non negative with $\tau = \max(\tau_{1J})(i,j=1,2$.as is shown in the appendix, the following system dominates (6.1*)*

$$\left.\begin{aligned}\frac{dZ_1(t)}{dt} &= -m_{11}\,Z_1(t) + b_{11}Z_1(t-\tau_{11}) - m_{12}Z_2(t-\tau_{12}),\\ \frac{dZ_2(t)}{dt} &= -m_{22}\,Z_1(t) + m_{21}Z_1(t-\tau_{21}) - b_{22}Z_2(t-\tau_{22}),\end{aligned}\right\} \quad t>0, \quad (6.2)$$

$Z_i(s) = |Y_i(s)|, \qquad i = 1,2 \;\; for\; s \le 0 \;\; for\; which\; Y_i(s)\; is\; defined.$

Thus the asymptotic stability of the trivial solution of (6.2) implies that of (6.1). it is also shown in the appendix that the solutions of (6.2) are nonnegative for all t $\geq$ 0. we will use the following vector-matrix notation:

$$\left.\begin{aligned}&Z = (Z_1, Z_2)^{\tau}, \qquad ||Z|| = |Z_1(t)| + |Z_2(t)| = Z_1(t) + Z_2(t)\\ &B = \begin{bmatrix} b_{11} & m_{12} \\ m_{21} & b_{12}\end{bmatrix}, \quad ||B|| = max \begin{matrix} b_{11} + & m_{21} \\ b_{21} + & m_{12}\end{matrix},\\ &A = \begin{bmatrix} -m_{11} & 0 \\ 0 & -m_{22}\end{bmatrix}.\end{aligned}\right\} \quad (6.3)$$

The logarithmic norm or the measure of any square matrix B, denoted by $\mu(B)$is defined as follows:

$$\mu(B) = \lim_{\theta \to 0+} \frac{/\!/I + \theta B/\!/ - 1}{\theta}$$

($I$ being an identity matrix of the same size as $B$).The matrix measure compatible with the norms in (6.3) is given by

$\mu(B) = \max_J [b_{jj} + \sum_{\substack{i=1 \\ i \neq J}}^{2} | b_{jJ} | ]$ $B=(b_{1J})$.

We refer to Coppel [11]or Vidyasager [31]for a number of properties of matrix measure We have shown in the appendix that the derivative of $H(Z_1,Z_2)(t)$(seeA.10)satisfies the relation

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq [-m_{11}+b_{11}+m_{21}]Z_1(t)+[-m_{22}+b_{22}+m_{12}]Z_2(t)$, and hence we have

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq \{\mu(A) + /\!/B/\!/\}\{ Z_1(t)+ Z_2(t)\}$

If we assume for the matrix norm $\|B\|$ as in (6.3)that

$$\mu(A) + /\!/B/\!/ < 0 \; or \; -\mu(A) > /\!/B/\!/$$

Then from the results of the appendix (Lemmas 5 and 6) it will follow that $Z_1(t)+Z_2(t) \to 0$ as $t \to \infty$ and thus we have the following:

Lemma 6.1 Assume that the Coefficients of the linear system (6.1)satisfy the condition (6.7). Then the trivial solution of(6.1)is asymptotically stable.

When (6.7)holds an estimate for the decay rate of solutions of (6.1)can be obtained as follows A necessary and sufficient condition for the asymptotic stability of the trivial solution is that all the roots of the characteristic equation

Det $G(\lambda,T_{iJ})=0$

Have negative real parts where

$$G(\lambda,T_{1J})=\begin{bmatrix} \frac{\lambda+m_{11}}{0} & \frac{0}{\lambda+22} \end{bmatrix} - \begin{matrix} b_{11}^{\,e-\lambda T_{11}} \\ \end{matrix} \begin{matrix} -m_{12}^{\,e-\lambda T_{12}} \\ -m_{21}^{\,e-\lambda T_{21}} \end{matrix} \quad \begin{matrix} \\ b_{22}^{\,e-\lambda T_{22}} \end{matrix}$$

$=\lambda I - A - B(\lambda,T_{iJ})$(say),

A positive number $\alpha$(if it exists)is said to be a decay rate of solution of (6.1) if and only if every root of (6.8)satisfies the condition

$Re(\lambda)+\alpha < 0$.

When (6.11)holds the solution of (6.1)will be of the order of $e^{-at}$ and hence the supremum of all such $\alpha > 0$ $a$ satisfying (6.11)will be a desirable parameter to estimate However such a supremum of all such (6.11)may not be possible to determine since it has to be done directly from (6.8)-(6,10)(see also Section 3).We will obtain an estimate of $\alpha$ (not necessarily the supremum)I (6.11)by the following procedure.

Replace λ byz-$\mathfrak{a}$ in (6.8)and define

$G(z-\mathfrak{a},T_{IJ})=M(z)=(z-a)I-A-B(z-\mathfrak{a},T_{IJ})$.

By Lemma 6.1 a Sufficient condition for all roots of M(z)=0 to have negative real parts equivalently to satisfy

$\mathrm{Re}(\lambda)< -a$

is the following :

$-\mu(aI+A) > \|B_1(a,T_{ij})\|$,

Where

$B_1(a,T_{ij})=\begin{pmatrix} b_{11}e^{-\lambda T_{11}} & -m_{12}e^{-\lambda T_{12}} \\ -m_{21}e^{-\lambda T_{21}} & b_{22}e^{-\lambda T_{22}} \end{pmatrix}$

Since $\mu(\mathfrak{a}\, I+A)=\mathfrak{a}+\boldsymbol{\mu}(\boldsymbol{A})$ we have from (6.13) that

$\mathfrak{a}+\mu(A)+\|B_I(\mathfrak{a},T_{IJ})\|< 0$

Is a sufficient condition for Re $(\lambda) < -a$If we denote the left side of (6.15)by $\mathfrak{a}f(\mathfrak{a},T_{Ij})$ we note that

$f(0,T_{Ij})=\|B\|+\mu(A)<0$ $if$ (6.7)$holds$;

Al so $\partial f/\partial\mathfrak{a}>0, f$ is continuous and monotonic in $(\mathfrak{a})$;$\to\infty as\ \mathfrak{a}\to\infty$; thus by the continuity of $f$ on $\mathfrak{a}\geq 0$, there exists an interval of the from $[0,^{*})$such that

$f(0,T_{Ij})< 0$ $for\ \mathfrak{a}\in[0,\mathfrak{a}^{*})$

And hence

$f(\mathfrak{a},T_{Ij})=0$

Thus an upper bound $\mathfrak{a}^{*}$ of $\mathfrak{a}$ for which we have $f(\mathfrak{a},T_{Ij})< 0$ $\mathfrak{a}\in[0,\mathfrak{a}^{*})$ can be obtained as the unique (by the monotonicity of $f$ on $\mathfrak{a}$) positive root of$\max\left\{\begin{matrix} b_{11}e^{-\lambda T_{11}}-m_{12}e^{-\lambda T_{21}}+\mathfrak{a}-\min\ m_{11}=0 \\ b_{22}e^{-\lambda T_{22}}-m_{21}e^{-\lambda T_{12}} \end{matrix}\right\}$ $m_{22}$

Where the maximum in (6.18) is taken over$b_{11}, b_{22}, m_{12}, m_{21}$ and $T_{ij}\,(i,=j\ 1,2)$. A numerical procedure can be devised to solve equations of the form (6.18) in concrete cases. If one is interested in a smaller bound on the decay rate it can be obtained as the unique root of

$\|B\| e^{ar} + a+\mu(A)= 0, \tau = \max_{i,J} T_{ij}$ . (6.19)

It will follow from (6.19) than an increase in T

Will reduce the value of the root of (6.19). We remark that since our estimate $a^*$ of dacay rate is based on a sufficient condition of the form (6.7), we can conclude that the solotions of (6.1) will converge at least as fast as $e^{-a*t}$ where $a^*$ is the unique root of either of (6.18) or (6.19).

**Implications**

Authors have formulated a class of two systems model systems staring from the elementary principal that the rate of change of a population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) equals the creation rate minus the destruction rate. This is a departure from the conventional model building in population ecology where one usually starts from the Malthusian formulation, now known as the Kolmogorov formulation in the nonlinear context. In our analysis of the models proposed we have obtained sufficient conditions for the global existence of nonnegative solutions of the model systems (this is trivial in the Kolmogorov formulation) and the other desirable properties such as the existence of suitable steady states. The two classes of models proposed for competition and debt contain intraspecific self – regulation biologists that each systems is under the action of a self-regulating negative feedback force, opinions and arguments differ regarding the mechanism by which such an effect is realized; we refer for more details on this aspect to a collection of articles compiled by Tamarin [30].

Gopalasamy's main result is the following : if a self –regulating negative feedback mechanism (intraspecific oppugn ant) can act without time delay and if such an effect can keep an ecosystem ( a community of one or more species of populations) asymptotically stable ( at least locally about a non trivial steady state, then time delays of any magnitude in any other interaction involving positive or negative feedbacks cannot destabilize such a stable ecosystem ( in a mathematical sense); that is such delays cannot cause a hopf-type bifurcation leading to persistent and undamped fluctuations. The delay-induced bifurcations to oscillations in models have been investigated by many authors. (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Qualitatively result is not new to the literature the effects of delays on otherwise stable and unstable systems have been considered by Cushing[12]where (pages 35-39)he has shown that an otherwise unstable system can be made less unstable by delays, Such an effect of delays on unstable systems appear to have been first proposed by Beddinton and May [2]who have also shows that delays can make stable systems less stable leading to a conclusion "time delays are not necessarily destabilizing" It is also shown by Cushing([12], pages 75-103)that while time delays cannot stabilize(in the mathematical sense)an otherwise unstable two species system in competition or mutualism time delays in there species systems can stabilize otherwise unstable systems (even in a mathematical sense)Cushing and Saleem [14]have formulated a model where delays in the positive feedback can have a stabilizing effect: Hale[18]discusses briefly the stabilizing nature of delays and given an example to show that an increase in delays can make an unstable system stable Thus for a better understanding of the effects of delays on otherwise stable systems one has to estimate the rate of decay of perturbations (or the characteristic return time )in terms of the parameters of the system especially when the delays do not cause a mathematical instability This aspect has been investigated by a procedure slightly different from that of the characteristic equation we have proposed (section6)

method for the estimation of decay rate of perturbations when there are multiple delays Since many species systems with multiple delays cannot be analysed (due to complexity)by the characteristic equation our methods using comparison lead to easily verifiable sufficient conditions The sufficient conditions obtained are in the form of a diagonal dominance type; these conditions can be weakened to quasi-diagonal dominance type as in Ladde [20], Siljak[20]Anderson [1], Lewis and Anderson [21,22](Ladde's result [20] is delay dependent due to the fact that delay appears in self –regulation also).Verification of quasi –diagonal dominance conditions in arbitrary matrices (except for Metzler matrices )is an impossible task and quasi –diagonal dominance is hard to interpret.

If follows from our discussion that those confederation systems which display delay induced instability leading to permanent and undamped oscillations should have one or more of the following aspects:

(1)time delays are in intraspecific interaction terms (and possibly in other terms also): (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

(2) Intraspecific self-regulation terms are absent in density dependent from in one or more populations of the system;

(3) Intraspecific negative feedback effect are not strong enough to maintain the stability of the relevant steady state of the nondelayed system.

We can now infer how the delay induced stability leads to fluctuations; a primary mechanism of such fluctuations can be described at least theoretically as follows : in its attraction of a system to a steady state, intraspecific self-regulation fails to be effective and the system overhoots the steady state and at such higher densities intraspecific competition becomes operationally effective leading the system to understood the steady state and this process is repeated again and again resulting in undamped oscillations about the steady state. It follows from our results that if intraspecific competition is sufficiently strong, the system cannot become oscillatory even when there are time delays, so long as they do not interfere with the strength of the intraspecific self – regulation.

We are thus led to the question of existence and importance of intraspecific competition in population systems. In the introduction to an anthology of articles by several authors, Tamarin [ 30] begins with the statement "all animal populations have the reproductive capacity to increase without limit, yet none di" and in his concluding remarks writes " this leaves us with only one great concept in the area of confederation do main regulation: that all credit-debit are regulated" More specifically this is applicable to all the confederation of (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

In his discussion on the effect of competition on the densities of populations, Nicholson ([ 25], $page$ [ 147] ) has reached the conclution "Intraspecific competition automatically regulates the severity of its action to the requirements of each denizen , provided the inherent resistance of the environment is sufficiently low to permit the species to exist." In order to look for a mechanism by which intraspecific competition plays its role of decreasing growth rate as confederation density increases, Calhoun [ 7] set up an experiment with Norway rates (Rattus norvegicus) and found that increased density leads to increased hostility which disturbs the psychology and physiology causing decreased reproductive output. Many other mechanisms by which intraspecific compition self-regulate populations have been put forward by a number of authors (see Tamarin [30]). This is true of

(confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

If we accept the view that intraspecific competition is really predominant in population regulation, then we can understand how and why time delays destabilize stable system leading to oscillations; time delays render intraspecific competitive regulation territorially (in case of particles we take the field in which it operates) weak so that self –regulatory forces become temporarily insufficient. If we now agree that all population models should incorporate self-regulating intraspecific competition terms in surfeit strength along with other tropic intraspecific competition terms in sufficient strength along with other atrophic interactions then it is not fair to say " time delays destabilize" and we note from our analysis that it is usually a failure in self- regulation which can cause instability rather than delays by themselves.

We will give one example to illustrate our point. Consider a credit-debit system in which the debits is self –regulating by a logistic load as follows: (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

$$\frac{dx}{dt} = r\left(x - \frac{x^2}{K}\right) - axy = b_I(x) - m_I(x,y) \quad (say),$$

$$\frac{dy}{dt} = \beta xy - \delta y = b_2(x,y) - m_2(y) - \quad (say),$$

Where $K, a, \beta, \delta$ are positive constants such that the steady state $x^* \; y^*$ exists where

$$x^* = \frac{\delta}{\beta}; \quad y^* = \frac{r}{a}\left(-\frac{\delta}{K\beta}\right) > 0.$$

In this example the predator self –regulation $(\partial m_2/\partial y)$ is not sufficiently strong in the sense

$$x^* = \left(\frac{\partial m_2}{\partial y}\right)_{(x^* y^*)} \succ \left(\frac{\partial b_2}{\partial y} + \frac{\partial m_1}{\partial y}\right)_{(x^* y^*)}$$

If in such a model there is a time delay in the credit creation rate, such a delay can lead to instability of $(x^* \; y^*)$ of setting up undamped oscillations. If however we modify the growth rate by the inclusion of a density dependent self-regulation so that credit

$$\frac{dy}{dt} = \beta\, xy - \delta y - \eta y^2$$

For some positive constant η then it is possible, depending on η, for such a modified system to remain stable without oscillations even if there is a time delay in the credit creations rate. We have another completely different justification to have $\eta y^2$ in the credit dynamics: for instance if the debit is removed from this cannot happen. While this example serves to illustrate an aspect of our result, our advocacy of the significance is not based on this example.

Gopalasamy has thus derived conclusion that "time delays need not destabilize" in models with several discrete delays. However we have provided in the appendix the necessary tools (Lemmas 3,4,5) to prove all our results when the relevant delays are all (or some) continuously distributed. Some authors (Caswell [9], May [23], Cushing [12] appear to regard continuously distributed dalays as financially/cosmologically credit wise obtained when he subjected algae to a variable nitrate environment and obtained a better fit for his data with a continuous delay. The questions of whether a discrete or continuous delay is more realistic may well depend on what process is being modeled. For instance an increased supply of some resource at some point in time may trigger reproduction at a particular moment in time later irrespective of the conditions of the resource during the intervening period. If the duration of increased resource supply necessary to elicit such a response is short compared to the intervening period before the reproduction occurs, then a discrete time delay may be appropriate. In a large number of cases of animal. Note fungus algae are prey. Here it shall be an account which loses the energy, amount, or matter for that matter. Same result holds good for (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Populations of particles, atoms, molecules, propositions, systems types chosen the past dietary and nutritional history of an animal over a long period plays an influential role in determining the current bahaviour of the dietary and nutritional history of an animal over a long period plays an influential role in determining the current behavior of the animal and in such cases continuous time delays will be appropriate.

In conclusion the following comments are appropriate. The time delays since the current creation rate is usually due to a contribution of those mature animals living a few units of time earlier depending on the type of animal or organism; the time delays appearing in interspecific interactions can be due to the fact that in order to successfully interact and compete interferentially a developmental maturity may be necessary. In the case of predation, delays in the birth rate of prey can be due to delayed maturation while that in the birth rate of a predator is due to gestation delay with delayed fecundity; the delays in the predation rate are again are again as before due to the maturity needed for a predator to successfully organize and succeed in an attack.

In our analysis we have assumed throughout a temporally uniform environment neglecting dispersal of populations and their age structure. We expect our results to hold under suitable conditions both in temporally periodic environments with the steady state of a constant environment being replaced by a periodic solution (nonstationary steady state) and when the effects of dispersion and age structures are considered. It is worthwhile to modify our purely ecological, cosmological, financial, propositional, problematic systems models so as to provide them with a genetic foundation and examine the effects of time delays since forces of selection act with some time delay. Such a generalization will definitely lead to a better understanding of "how confederations regulate themselves".

## Appendix

First we derive a comparison result.

Lemma 1. Consider the delay differential system

$$\left.\begin{aligned}\frac{dY_1(t)}{dt} &= -m_{11}\ Y_1(t) + b_{11}Y_1(t-\tau_{11}) - m_{12}\ Y_2(t-\tau_{12}),\\ \frac{dY_2(t)}{dt} &= -m_{22}Y_2(t) - m_{21}Y_1(t-\tau_{21}) + b_{22}\ Y_2(t-\tau_{22}),\end{aligned}\right\} (A.1),$$

Where the $m_{ij}$, $b_{ij}$ $(i.j = 1,2)$ are positive while $\tau_{ij}$ $(i.j = 1,2)$ nonnegative with $\tau = \max \tau_{ij}$, $i,j = 1,2$. Then

$$\left.\begin{aligned}\frac{d|Y_1(t)|}{dt} &\le -m_{11}\,|Y_1(t)| + b_{11}|Y_1(t-\tau_{11})| + m_{12}\ Y_2(t-\tau_{12})|,\\ \frac{d|Y_1(t)|}{dt} &\le -m_{22}\,|Y_2(t)| + m_{21}|Y_1(t-\tau_{21})| + b_{22}\ |Y_2(t-\tau_{21})|,\end{aligned}\right\} (A.2)$$

Where

$$\frac{d}{dt}\,|\,Y_1(t)\,| = \left\{\begin{aligned}&\frac{dY_1(t)}{dt}\ if\ \ Y_1(t) \ge 0,\\ &\frac{dY_1(t)}{dt}\ if\ \ Y_1(t) < 0,\end{aligned}\right\}\ i = 1,2. \qquad (A.3)$$

The proof is by direct verification.

LEMMA 2. Consider now (A.2) and the following system

$$\left.\begin{aligned}\frac{dZ_1(t)}{dt} &= -m_{11}\ Z_1(t) + b_{11}Z_1(t-\tau_{11}) + m_{12}\ Z_2(t-\tau_{12}),\\ \frac{dZ_2(t)}{dt} &= -m_{22}Z_2(t) + m_{21}Z_1(t-\tau_{21}) + b_{22}\ Z_2(t-\tau_{22}),\end{aligned}\right\} \qquad (A.4)$$

Together with the initial conditions

$$Z_1(s) > |Y_1(s)|,\ \ s \in [-\tau, 0],\ i = 1,2. \qquad (A.5)$$

Then

$$|Y_1(t)| < Z_1(t) \text{ for } t \ge 0 \quad \text{and} \quad i = 1,2. \qquad (A.6)$$

PROOF. Let $U_1(t) = Z_i(t) - |Y_1(s)|$, $i = 1,2$. $t \ge -\tau$. Then we have from (A.2) and (A.4),

$$\left.\begin{aligned}\frac{dU_1(t)}{dt} &\ge -m_{11}\ U_1(t) + b_{11}U_1(t-\tau_{11}) + m_{12}\ U_2(t-\tau_{12}),\\ \frac{dU_2(t)}{dt} &\ge -m_{22}U_2(t) + m_{21}U_1(t-\tau_{21}) + b_{22}\ U_2(t-\tau_{22}),\end{aligned}\right\} t > 0, \quad (A.7)$$

With

$U_1(s) = Z_i(s) - |Y_i(s)| > 0, \quad i = 1,2, s \in [-\tau, 0]$. (A.8)

Consider (A.7)-(A.8) first on the interval $[\tau, 0]$ : It will follow from (A.7) that $U_1(t) \geq U_1(0) \exp [-m_{11} t]$

$+\int_0^t [b_{11} U_1(s-\tau_{11}) + m_{12} U_2(s-\tau_{12})] \exp [-m_{11}(t-s)ds > 0,]$ (A.9)

For all $t \in [0, \tau^*]$ since $U_i(s) > 0$ on $[-\tau^* 0]$ where $\tau^*$ = positive minimum of $(\tau_{11,}\tau_{12,}\tau_{21,}\tau_{22,})$. Similarly $U_2(t) > 0$ for $t \in [0, \tau^*]$. Now by a similar argument it will follow that $U_1(t) > 0, U_2(t) > 0$ for $t \in [\tau^*, 2\tau^*]$. Repeating the above procedure we obtain the result (A.6).

It follows from Lemmas 1 and 2 that in order to prove theorem 3.1 it is enough to prove that $Z_i(t) \to 0$ as $t \to \infty$ for $i = 1,2$ . which we will do as follows.

PROOF OF THEOREM 3.1 Define a Lyapunov-like function

$H(Z_1,Z_2)(t)=Z_1(t)+Z_2(t)$

$+\int_{1-T}^t b_{11} Z_1(s)ds + \int_{1-T}^t m_{12} Z_2(s)ds$

$+\int_{1-T}^t m_{21} Z_1(s)ds + \int_{1-T}^t b_{22} Z_{22}(s)ds$

Since each term of H in (A.10)is nonnegative for t>0, if we show that $H \to 0$ as $t \to \infty$ it will follow from Lemma 2 that

$0 \leq |Y_{1,}(t)| \leq Z_1(t) \to 0$ as $t \to \infty$,

Implying that the perturbations $Y_{1,}$, $Y_2$, governed by (A.1)approach zero as $t \to \infty$, and this will prove our result Calculate the rate of change of H along the solutions of (A.4)so that

$\frac{dH}{dt} = \frac{\partial H}{\partial t} + \frac{\partial H}{\partial Z_1}\frac{dZ_1}{dt} + \frac{\partial H}{\partial Z_2}\frac{dZ_2}{dt}$

$=-[Z_1(t)\{m_{11} - (b_{11} + m_{12})\} + Z_2(t)\{m_{22} - (b_{22} m_{12})\}] < 0$

On using (2.10)since $Z_1(t),Z_2(t)$ are positive for all t>0. Thus $H(Z_1, Z_2)(t) \to 0$ as $t \to \infty$ since H is bounded below by zero: it follows from Lemma 6 below that $Z_1(t) \to 0, Z_2(t) \to 0$ as $t \to \infty$ the proof is complete.

The following set of results are intended to be a guide to prove result similar to the above in the case of continuously distributed delays An analogue of Lemma I is the following.

LEMMA 3 Let $K_{ij}(s)$,$s \in [0,\infty)$, $i, j = 1,2$,denote nonnegative continuous functions integrable on $[0,\infty)$ normalised such that

$\int_0^\infty s k_{ij}(s)ds < \infty$; $\int_0^\infty k_{ij}(s)ds = 1$, I,j=1,2.

Consider the intergrodifferential system

$\frac{dY_1(t)}{dt} = -m_{11}Y_1(\mathrm{t}) + b_{11}\int_0^\infty k_{\mathrm{ij}}(\mathrm{s})\, Y_1(\mathrm{t}-\mathrm{s})\mathrm{ds}$

$-m_{12}\int_0^\infty k_{12}(\mathrm{s})\, Y_2(\mathrm{t}-\mathrm{s})\mathrm{ds}$

$\frac{dY_2(t)}{dt} = -m_{22}Y_2(\mathrm{t}) + m_{21}\int_0^\infty k_{21}(\mathrm{s})\, Y_1(\mathrm{t}-\mathrm{s})\mathrm{ds}$

$+b_{22}\int_0^\infty k_{22}(\mathrm{s})\, Y_2(\mathrm{t}-\mathrm{s})\mathrm{ds}$

*Where* $m_{ij}, b_{ij}\ (i,j=1,2)$ *are positive constants Then*

$$\frac{d|Y_1(t)|}{dt} \leq -m_{11}|\,Y_1\,(t)\,| + b_{11}\int_0^\infty k_{11}\,(s)|Y_1(t-s)|ds,$$

$$+m_{12}\int_0^\infty k_{12}\,(s)|\,Y_2(t-s)|ds,$$

$$\frac{d|Y_2(t)|}{dt} \leq -m_{22}|\,Y_2\,(t)\,| + m_{21}\int_0^\infty k_{21}\,(s)|Y_1(t-s)|ds,$$

$$+b_{22}\int_0^\infty k_{22}\,(s)|\,Y_2(t-s)|ds,$$

*Where* $d|Y_1(t)||/dt$ *is defined by* (A.3).

Proof is again by direct verification.

The proof of the following lemma is similar to that of Lemma 2 above.

Lemma 4. *Consider (A.14) and the following intergrodifferential system:*

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1\,(t) + b_{11}\int_0^\infty k_{11}\,(s)Z_1(t-s)ds$$

$$+m_{12}\int_0^\infty k_{12}(s)Z_2\,(t-s)ds,$$

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1\,(t) + b_{11}\int_0^\infty k_{11}\,(s)Z_1(t-s)ds$$

$$+b_{22}\int_0^\infty k_{22}\,(s)Z_2\,(t-s)ds,$$

Together with the initial conditions

$$\infty > Z_i\,(s) > |\,Y_i\,(s)|for\ s \in (-\infty, 0).$$

(A.16)

Then $|Y_i(t)| < Z_i(t)for\ t \geq 0, i = 1,2.$

The following lemma will be useful to prove a result analogous to that of Theorem 3.1 corresponding to a class of continuously distributed time delays.

LEMMA 5. Let $Z_1(t), Z_2\,(t)$ be the solutions of (A.15)- (A.16) where $(Y_1\ Y_2)$ corresponds to an arbitrary but fixed solutions of (A.13). Then as a consequence of Lemma & it will follow that

$$0 \leq |Y_i\,(t)|\ Z_i\,(t) \to \infty \leq as\ t \to \infty, \qquad i = 1,2,$$

(A.17)

Whenever

$\left\{\begin{matrix} M_{11} > b_{11}+m_{21}, \\ M_{22} > b_{22}+m_{12} \end{matrix}\right\}$

PROOF. Define a Lyapunov –like function H where

$H(Z_1,Z_2)(t)=Z_1(t)+Z_2(t)+b_{11}\int_0^\infty k_{11}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$

$+m_{12}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$

$+m_{21}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$

$+b_{22}\int_0^\infty k_{22}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$

Where $(Z_1,Z_2)$is any solution of (A.15)-(A.16) for arbitrary $(Y_1,Y_2)$defined on $(-\infty, 0]$ Computing the rate of change of H in (A.19)along the solution of (A.15)we derive that (on using (A.12))

$\left(\frac{dH(Z_1Z_2(t)}{dt} = [Z_1(t)\{m_{11} - (b_{11} + m_{21})\} + Z_2(t)\{m_{22} - (b_{22} + m_{12})\}]\right.$

$< 0$(using (A.18)) for $Z_1(t) > 0$ $and$ $Z_2(t) > 0$.

By Lemma 6 below the result will follow

Lemma 6. Let H(t)=H$(Z_1, Z_2)(t)$ be as defined in (A.10) Suppose

$$m_{11} > b_{11} + m_{21} \text{ and } m_{22} > b_{22} \mid m_{12}$$

Then

$$Z_1(t) + Z_2(t) \to 0 \text{ast} \to \infty.$$

PROOF. Suppose (A.21) dose not hold; then there exists at least one sequence

$\{t_n\} \to \infty$ as $n \to \infty$ such that

$$[m_{11} - (b_{11} + m_{21})]\, Z_1(t_n) + [m_{22} - (b_{22} + m_{12})]Z_2(t_n) > \varepsilon \qquad (A.22)$$

For some sufficiently small positive number $\varepsilon$ and hence

$$\frac{dH(t_n)}{dt} < -\varepsilon, n = 1,2,3, \ldots.$$

By the continuity and differentiability of H with respect to t,

$H(t_n - \varepsilon) - H(t_n) = \frac{dH(t^*)}{dt}(-\varepsilon), \quad t_n - \varepsilon < t^* < t_n$

If $\varepsilon$ is small we will have

$$= \frac{dH(t^*)}{dt}\varepsilon/2$$

So that

$$H(t_n) \leq H(t_n - \varepsilon) - \varepsilon^2/2$$

$$\leq H(t_{n-1}) - \varepsilon^2/2$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2)$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2):$$

$$\leq H(t_0) - n(\varepsilon^2/2): t_0=0.$$

(A.12) and Lemma 5 will imply that $H(t_0) < \infty$ for any $t_0 > 0$. If n is sufficiently large it will follow that $H(t_n) < 0$ since $H(t_0) < \infty$,which contradicts the nonnegativity of H for all $t \geq 0$. Hence $Z_1(t) + Z_2(t) \to 0$ as $t \to \infty$ implying the asymptotic stability of the trivial solution of (A.1)(since $Z_1$ dominates $y_1$)

We note the result of Lemma 5 and 6 will remain valid if (A.18) is replaced by the condition $\mu(A) + /\!/B/\!/ < 0$

Where A and B are defined as in (6.3).

We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems

**Gopalsamy**

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# Model For Innocuously Benign Delay In A Periodic System

**Abstract**

Sufficient conditions are obtained for the existence of a unique linearly denomination stable positive periodic solution of a system model of two competitions in a periodic environment with time lags in interspecific interactions. It is shown that if the self- regulating intraspecific interaction effects are strong enough and act without time delays then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a periodic domination. An example in finance will make things clear. Corporate finance for one big plant generates lot of accounts of the workers whose salary shall be disbursed to the debit of cash credit account. Here Cash credit account is a resource as also the one to which debits are raised, a la predator. Some of the accounts from which the credits come are demand bills purchase, usance bills discount, vostro and nostro accounts in respect of foreign bills and Letters of credit. In economics inventories are used to produce output, output is sold to produce income which might be saved in part and recycled in investment. Cosmologically take an example of a Black hole which gormandizes matter and radiates energy. Notwithstanding information paradox, one thing is clear that matter is being created in to radiation. We show how the model could be applied to axiomatic predications, postulation alcovishness, phenomenological methodologies, propositional statement, and other problems in cosmology, philosophy, and even in Green House effects. We study the sine qua non properties of such a system. In fact these models can be applied to such unsolved problems like Yang mills Theory and mass gap, and other problems. As we apply the model, we shall go on dilating the position status and the applicability of the model. We shall not touch the electrical engineering problems which might rise questions like what is charge of a Bank. We may say when a customer enters a Bank; we can say the Bank is charged. These factors are still relativistic and we shall make an in-depth study as we progress. Gopalsamy's research was published in Australian Journal of Mathematical Society vide: B 25349=365 It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 4. Introduction

At the outset we assume that the system satisfies Haimovici is caution. The purpose of this article is to establish a set of sufficient conditions for the existence of a unique linearly asymptotically stable positive periodic solution for a nonautonomous system of the form

$$\frac{dx(t)}{dt} = x(t)f_1\big(t, x(t), y(t-\tau_1)\big),$$

**(1.1)**

$$\frac{dy(t)}{dt} = y(t)f_2\,\big(t, x(t-\tau_2), y(t)\big).$$

In the context of confederation dynamics a model system of the form (1.1) means that the interspecific interactions occur with time delays $\tau_1 \geq 0, \tau_2 \geq 0$ in a periodic environment, if $f_1$ *and* $f_2$ are assumed to be periodic in $t$ with a common period $\omega > 0$. Note the interactions are inter specific.

Although model systems in periodic domains have been investigated by a large number of authors (for some references see Gopalsamy (4], [5]), periodically varying environments in time delayed model systems of two or more interacting species have not been investigated. The need for Multitype system models in periodic environments and time delayed interactions are evident from the extensive literature of that type relating to single species. While it is known that periodic ecosystems can support more species than uniform environments (Gopalsamy [4]), time delays might destabilize the system leading to a possible loss of one or more species. Furthermore, periodic environments in population dynamics have some relevance to natural selection and evolution; we will consider this aspect briefly at the end.

The principal finding of the following analysis of (1.1) is that, under suitable hypotheses on $f_1$ *and* $f_2$, the system (1.1) has a linearly asymptotically stable unique positive periodic solution for all nonnegative delays $\tau_1, \tau_2$ (while our numerical experiments suggest that the periodic solution of (1.1) is globally asymptotically stable and is independent of the delay parameters $\tau_1, \tau_2$ it is not known how one can prove such a strong result). Since time delays are usually associated with possible destabilization of otherwise stable systems, we can describe the time delays appearing on (1.1) as "harmless".

It has been recently shown by the author (Gopalsamy [6]) that if the self –regulating negative feedback in each denomination is strong enough and acts without time delays, then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a temporally uniform domain. The result of this paper is thus an extension to a periodic environment of the "harmless" nature of interspecific interaction delays in a temporally constant environment mentioned above.

## 5. Existence and uniqueness of periodic solutions

We will now formulate the problem of existence of periodic solutions of the periodic delay differential system

$$\frac{dx_1(t)}{dt} = x_1(t) f_1\big(t, x_1(t), x_2(t - \tau_{12})\big),$$

$t > 0,$

**(2.1)**

$$\frac{dx_2(t)}{dt} = x_2(t) f_2\big(t, x_1(t-\tau_{21}), x_2(t)\big),$$

(Where $\tau_{12}$ , $\tau_{21}$ are nonnegative constants) in a form suitable for the application of fixed point theory of positive operators. We make the following assumptions regarding $f_1$ and $f_2$ in (2.1).

($H_1$) The partial (Frechet) derivatives of $f_1, f_2$ satisfy the following:

$$\frac{\partial f_i}{\partial x_j} < 0, \qquad i,j = 1,2 \text{ For all } t > -\tau \text{ where } \tau = \max(\tau_{12}, \tau_{21}). \qquad \textbf{(2.2)}$$

($H_2$) $f_1(t, x_1, x_2) \equiv f_1(t + \omega, x_1, x_2)$ for all $t > -\tau, x_1 \geq 0, x_2 \geq 0.$ **(2.3)**

($H_3$) There exist positive constants $b_{il}$, $b_i^u$, $a_{il}$, $a_i^u$ ( $l = 1,2$) such that

$$b_{il} - a_{il} x_i < f_i(t, x_1, x_2) < b_i^u - a_i^u x_i,$$

$$b_{il}/a_{il} < b_i^u / a_i^u, \qquad \textbf{(2.4)}$$

For $i$= 1, 2 and $(x_1, x_2) \in R_+ \times R_+$. If $\tau = 0$ then one can show that (2.1)-(2.4) will imply that the nondelayed periodic system (2.1) will have an asymptotically stable strictly positive periodic solution (for details see Gopalsamy [5]); this result will become a special case of what we obtain in the following.

We will use the following notation:

$$B_\omega^\tau = \{h: [-\tau, \infty) \to R, h \text{ is continuous}, h(t + \omega) = h(t), t \geq -\tau\},$$

$$B = B_\omega^\tau \times B_\omega^\tau,$$

$$K_1 = \{h \in B_\omega^\tau | h(t) \geq 0, t \in [-\tau, \infty)\},$$

$$K = K_1 \times K_1,$$

$$\|h\| = \left\{\sup_{t \in [0,\omega],} |h(t)|, h \in B_\omega^\tau\right\}$$

$$\|x_1, x_2\| = \max\{\|x_1\|, \|x_2\|, (x_1, x_2) \in B\},$$

$$G_1 = \{(x_1, x_2) \in K, \|(x_1, x_2)\| \geq b_l = \min(b_{1l}/a_{1l}, b_{2l}/a_{2l})\}$$

$$G_2 = \left\{(x_1, x_2) \in K, \|(x_1, x_2)\| \geq b^u = \max\left(\frac{b_1^u}{a_1^u}, \frac{b_2^u}{a_2^u}\right)\right\},$$

$\bar{G}$ = closure of $G_2$ *in* $B$ and

$\frac{\bar{G}_2}{G_1}$ = complement of $G_1 in\ \bar{G}_2;\ G_1 \subset G_2$ by **(2.4)**

The cone $K = K_1 \times K_1$ induces a partial order in $B$ and we will write $(x_1, x_2) \leq (y_1, y_2)$ if and only if $x_1 - y_1 \in K_1$. we define an operator $N$ on $B$ as follows

$$N(x_1, x_2)(t) = \begin{cases} N_1(t, x_1, x_2) \\ N_2(t, x_1, x_2) \end{cases} \quad \text{for t> } 0, \qquad \textbf{(2.5)}$$

Where

$$N_1(t, x_1, x_2) = \int_0^{\omega} G_{11}(m, \omega, t-s) g_1\left(m, s, x_1(x) x_2(s-\tau_{12})\right) ds$$

$$+ \int_0^1 G_{12}(m, 0, t-s) g_1\left(m, s, x_1(s), x_2(s-\tau_{12})\right) ds,$$

$$N_2(t, x_1, x_2) = \int_0^{\omega} G_{22}(m, \omega, t-s) g_2\left(m, s, x_1(s-\tau_{21}), x_2(s)\right) ds \qquad \textbf{(2.6)}$$

$$+ \int_0^1 G_{21}(m, 0, t-s) g_2\left(m, s, x_1(s-\tau_{21}), x_2(s)\right) ds,$$

For $t > 0$, with

$$g_1\left(m, s, x_1(s), x_2(s-\tau_{12})\right) = m x_1(s) + x_1(s) f_1\left(s, x_1(s), x_2(s-\tau_{12})\right),$$

**(2.7)**

$$g_2\left(m, s, x_1(s-\tau_{21}), x_2(s)\right) = m x_2(s) + x_2(s) f_2\left(s, x_1(s-\tau_{21}) x_2, (s)\right),$$

$$G_{ll}(m, \omega, t-s) = e^{-m\omega}(1 - e^{-m\omega})^{-1} e^{-m(t-s)}, \qquad i = 1, 2,$$

**(2.8)**

$$G_{ij}(m, 0, t-s) = e^{-m(t-s)}, i \neq j, i, j = 1, 2,$$

$$N(x_1, x_2)(t) = \begin{cases} N_1(n\omega + t, x_1, x_2) \\ N_2(n\omega + t, x_1, x_2) \end{cases} for\ t \in [-\tau, 0], \qquad \textbf{(2.9)}$$

*m* being a suitably selected positive constant (more will be said below) while *n* is a positive constant such that $n\omega > \tau$.

We note immediately the following from the definition of $N$: $if\ (x_1, x_2) \in \bar{G}/G_1$ then $N(x_1, x_2)(t)$ 0 and if $N(x_1, x_2)$ is periodic $\omega$ for $t$ 0 then $N\ (x_1, x_2)$ is periodic in *t* with period $\omega$ for $t$ −$\tau$. The existence of at least one positive periodic solution of (2.1)-(2.4) is established by means of a few lemmas and a result of Schmitt [9].

Lemma 1. *A necessary and sufficient condition for* $(x_1(t), x_2(t))$ , $t$ −$\tau$ *to be a periodic solution of (2.1)-(2.4) of period* $\omega$ *is that* $(x_1, x_2)$ *be a fixed point of N such that*

$$N(x_1, x_2)(t) = \big(x_1(t), x_2(t)\big) \quad and\ (x_1, x_2)\ \in K\ \cap \left(\frac{\bar{G}}{G_1}\right).$$

PROOF. The details of proof are similar to those of Lemma 2 of Busenberg and Cooke [1]. We first verify the continuity of $N(x_1, x_2)$ assuming that $N(x_1, x_2)(t)$ is periodic in *t* of period $\omega$. The continuity of $N(x_1, x_2)$ for $t > 0$ follows from the definition in (2.5)-(2.8) WHICH ALSO PROVIDES $\lim_{t\to 0+} N(x_1, x_2)(\mathrm{t})$ . It will follow from (2.9) and the assumed periodicity of $N(x_1, x_2)\ that\ N(x_1, x_2)$ is continuous

For $t$ −$\tau$. Now we will verify the periodicity of $N(x_1, x_2)$ in $t$. A simple computation leads to

$$N_1(t + w, x_1, x_2) - N_1(t, x_1, x_2) = e^{-mt} F_1(t),$$

**(2.10)**

Where

$$F_1(t) = \int_0^{\omega} G_{11}\ (m, \omega, t - s\ ) g_1\big(m, s, x_1(s), x_2(s - \tau_{12})\big) ds$$

$$+ \int_0^{\omega} G_{12}(m, 0, t - s) g_1\ \big(m, s, x_1(s), x_2(s - \tau_{12})\big) ds$$

$$+ \int_{\omega}^{\omega + t} G_{12}(m, 0, \omega - s) g_1\ \big(m, s, x_1(s), x_2(s - \tau_{12})\big) ds$$

$$- \int_0^{\omega} G_{11}(m, \omega - s) g_1\ \big(m, s, x_1(s), x_2(s - \tau_{12})\big) ds$$

$$+ \int_0^t G_{12}(m, 0, -s) g_1 \left(m, s, x_1(s), x_2(s - \tau_{12})\right) ds$$

**(2.11)**

Which on simplification using the periodicity of $f_1$ shows that $F_1(t) \equiv 0$ whenever $x_1, x_2, \in B$, thus implying the periodicity of $N_1$ in $t$ of period$\omega$. A similar argument for $N_2$ proves that $(x_1, x_2) \in B \Rightarrow N(x_1, x_2) \in B$.

Now suppose $(x_1, x_2)$ is a fixed point of *N and let* $(x_1, x_2) \in K \cap (\bar{G}_2/G_1)$. Then it is easily seen that $(x_1, x_2)$ is a solution of (2.1) and since $(x_1, x_2) \in K \cap (\bar{G}_2/G_1)$ it will follow that $(x_1, x_2)$ is a periodic solution of (2.1).

Conversely, let $(x_1, x_2)$ be a periodic solution of period $\omega$ of (2.1). Then $(x_1, x_2)$ is necessarily a periodic solution of

$$\frac{dx_1(t)}{dt} + mx_1(t) = g_1\left(m, t, x_1(t), x_2(t - \tau_{12})\right)$$

**(2.12)**

$$\frac{dx_2(t)}{dt} + mx_2(t) = g_2\left(m, t, x_1(t - \tau_{21}), x_2(t)\right),$$

For any arbitrary real constant *m* and hence for any fixed constant $m > 0$. It follows from (2.12) that any periodic solution of (2.1) and hence of (2.12) is a fixed point of $N(x_1, x_2)$. We claim that such a fixed point is in $K \cap (\bar{G}_2/G_1)$. Suppose $(x_1, x_2) \notin K \cap (\bar{G}_2/G_1)$. the following possibilities can arise:

(iv) Either $x_1, x_2$ or both $x_1$ *and* $x_2$ may become negative for some value of *t* and it will follow from (2.1) that subsequent value of $x_1, x_2$ or $x_1$ *and* $x_2$ will be negative and hence from (2.4) and (2.1) their derivatives will be negative for all subsequent *t* , implying that $x_1, x_2$ or $x_1$ *and* $x_2$ cannot be periodic.

(v) While $(x_1, x_2) \in K \subset B$ there may be a value of *t*, say $t^*$, such that $x_1(t^*) > b_1^u/a_1^u$ or $x_2(t^*) > b_1^u/a_1^u$ and in such a case $x_1(t) or\ x_2(t)$ cannot be periodic since respectively $dx_1(t^*)/dt < 0\ or\ dx_2(t^*)/dt < 0$, which will imply that the values of $x_1(t^*)\ or\ x_2(t^*)$ cannot be repeated from below.

(vi) Similarly if there is a $t^*$ such that $0 < x_1(t^*) < b_{1l}/a_{1l}$ *or* $0 < x_2(t^*) < b_{2l}/a_{2l}$ then such solutions cannot be periodic since $dx_1(t^*)dt > 0$ or $dx_2(t^*)dt > 0$, which will mean that the values of $x_1(t^*)$ or $x_2(t^*)$ cannot be repeated from above.

Thus every periodic solution of (2.1)-(2.4) is a fixed point of *N* in $K \cap \left(\frac{G_2}{G_1}\right) \subset B$.

**LEMMA 2.** *Let N be as in (2.5)-(2.6). Suppose* $\lambda \in (0,1) and \ \lambda N(x_1, x_2) = (x_1, x_2)$ *for some* $((x_1, x_2) \in K - \{0\}$. *then there exist* $\alpha_1^* > 0, \alpha_2^* > 0$ *such that*

$$0 < x_1(t) \geq \alpha_1^* < b_1^u / a_1^u,$$

$$0 < x_2(t) \geq \alpha_2^* < b_2^u / a_2^u,$$

*for* $t \in [-\tau, \infty)$.

PROOF. If $\lambda N(x_1, x_2) = (x_1, x_2)$ then from the definition of N we get for all $t > 0$ and $(x_1, x_2) \in K - \{0\}$,

$$\frac{dx_1(t)}{dt} = \lambda x_1(t) f_2(t, x_1(t), x_2(t - \tau_{12})) + x_1(t) m(\lambda - 1)$$

**(2.13)**

$$\frac{dx_2(t)}{dt} = \lambda x_2(t) f_2(t, x_1(t - \tau_{21}), x_2(t)) + x_2(t) m(\lambda - 1)$$

For all $t \geq 0$ such that $x_1(t) \geq b_1^u/a_1^u$ *or* $x_2(t) \geq b_2^u/a_2^u$ we have $\frac{dx_1}{dt} < 0$ *or* $\frac{dx_2}{dt} < 0$; but these are impossible since $x_1, x_2$ are periodic in t. Thus by the continuity and periodicity of $x_1, x_2$ we have $x_1(t) \geq \alpha_1^* < b_1^u/a_1^u$, $x_2(t) \geq \alpha_2^* < b_2^u/a_2^u$ for all $t \geq 0$; by the definition and periodicity the same will follow for all $t \geq -\tau$. A similar argument holds when $x_1(t) < 0$ or $x_2(t) < 0$, noting that in such a case $dx_1/dt > 0$ or $dx_2/dt > 0$, again contradicting the periodicity of $(x_1, x_2)$ in *t*.

LEMMA 3. *The operator N maps K∩ $(G_2/G_1)$ in to K and is completely continuous*

PROOF. From the definition of N and (2.4) we derive that

$$N_1(t,x_1,x_2) = \int_0^{\omega} G_{11}\,(m,\omega,t-s)x_1(s)[m+b_{1l}-a_{1l}x_1(s)]ds$$

$$+\int_0^{\omega} G_{12}\,(m,0,t-s)x_1(s)[m+b_{1l}-a_{1l}x_1(s)]ds.$$

**(2.14)**

For $(x_1,x_2)\in K\cap(G_2/G_1)$, we have $0<x_1<b_1^u/a_1^u$ and hence we get $N_1(t,x_1,x_2)>0$ for all $t$ by a proper choice of the positive constant say $m_1$.

We make a similar choice say $m_2$ for *m* for the case $N_1(t,x_1,x_2)$ and then choose *m* = max $(m_1,m_2)$. Thus we have

$$N: K\cap(G_2/G_1)\to K \qquad for\ t,$$

And by the periodicity of N we have the same for $t-\tau$.

To prove the complete continuity of N we proceed as follows. We know that if $(x_1,x_2)\in B$ then $N(x_1,x_2)\in B$ and so we can restrict *t* to the set $[0,n\,\omega], n\omega>r$. We write $N=L\circ Q$ where *L* is the linear operator, $L: B\to B$, $L(u_1,u_2)(t)$

$$=\int_0^{\omega} G_{11}(m,\omega,t-s)\begin{bmatrix}u_1 & (s)\\ u_2 & (s)\end{bmatrix}ds,+\int_0^{t} G_{12}(m,0,t-s)\begin{bmatrix}u_1 & (s)\\ u_2 & (s)\end{bmatrix}ds$$

**(2.15)**

And Q is the nonlinear operator

$$Q(u_1,u_2)(t) = \begin{matrix} g_1(m,t,u_1(t),u(t-\tau_{12})) \\ g_2(m,t,u_1(t-\tau_{21}),u_2(t)) \end{matrix} for\ t>0,$$

**(2.16)**

$$Q(u_1,u_2)(t) = Q\,(u_1,u_2)(n\omega+t)\quad for\ t\in[-\tau,0].$$

It can be found that

$$||L(u_1,u_2)(t)|| \leq max\left\{\frac{m+b_1^u}{m},\frac{m+b_2^u}{m}\right\}||u_1,u_2||,$$

**(2.17)**

$$\left\|\frac{d}{dt}L(u_1,u_2)\right\| \leq max\left\{\frac{m+b_1^u}{m},\frac{m+b_2^u}{m}\right\}||u_1,u_2||.$$

The complete continuity of the linear map $L$ follows from (2.17). Now since Q is continuous and maps bounded sets into bounded sets it will follow that $L \circ Q$ is completely continuous.

To prove the existence of periodic solution of (2.1)-(2.4) we will use the following result of Schmitt [9].

**THEOREM (Schmitt).** *Let $H_1$ and $H_2$ be nonempty bounded open neighborhoods of the origin in a Banach space B with a positive cone K such that $H_1 \subset H_2$ and let N: $K \cap (H_2/H_1) \to K$ be completely continuous and which satisfies the following conditions:*

***(c)*** *There exists $(z_1,z_2) \in K \subset B, \|(z_1,z_2)\| = 1$ and a $\mu$ satisfying $\mu > sup\{\|(x_1,x_2) - N(y_1,y_2)\|, (x_1,x_2) \in K \cap H_1, (y_1,y_2) \in K \cap \partial H_2\}$ such that all solution $(x_1,x_2) \in K \cap (H_2/H_1)$ of $(x_1,x_2) = N\,(x_1,x_2) + \lambda(z_1,z_2) with\ 0<\lambda<\mu\ satisfy\ (x_1,x_2) \notin \partial H_1$.*

*(d) All solution $(x_1,x_2) \in K \cap (H_2/H_1)$ of $(x_1,x_2) = \lambda\, N(x_1,x_2), 0<\lambda<1\ satisfy\ (x_1,x_2) \notin \partial H_2$.*

*Then N has at least one fixed point in $K \cap (H_2/H_1)$.*

We can now present our existence result.

**THEOREM 1.** *Assume that* $f_1, f_2$, *of* (2.1) *satisfy* (2.2)-(2.4). *then there exists at least one positive periodic solution of* (2.1)-(2.4).

PROOF. By Lemma 1, it is enough to show that the operator $N: K \cap (G_2/G_1) \to K \subset B$ has at least one fixed point in $G_2/G_1$ and we will do this by applying the result of Schmitt [9]. First we let

$$H_2 = \{(x_1, x_2) \in B | \, ||x_1, x_2|| < \max(b_1^u/a_1^u, b_2^u/a_2^u)\} \equiv G_2,$$

$$H_1 = \{(x_1, x_2) \in B | \, ||x_1, x_2|| < \min(b_{1l}/a_{1l}, b_{1l}/a_{1l})\} \equiv G_1.$$

By Lemma 3, $N: K \cap (H_2/H_1) \to K$ is completely continuous. Choose $\mu = \infty$ and $(z_1, z_2) \equiv (1.1)$. Then all solution $(x_1, x_2) \in K \cap (H_2/H_1)$ of $(x_1, x_2) = N\,(x_1, x_2) + \lambda(z_1, z_2), 0 < \lambda < \infty$, satisfy (2.1) and hence by their periodicity $(x_1, x_2) \notin \partial H_1$. The condition (a) of Schmitt's theorem is satisfied since $\mu = \infty$ by choice and $(x_1, x_2)$ and $(z_1, z_2)$ are periodic. By Lemma 2 the condition (b) of Schmitt's theorem is satisfied. Thus there exists at least one fixed point of N in $K \cap (H_2/H_1)$ and hence there exists at least one positive periodic solution of (2.1)-(2.4) and the proof is complete.

We remark that so far no restriction on the magnitudes of the delays $\tau_{12}, \tau_{21}$ has been imposed except that they be nonnegative.

## 6. Uniqueness of the periodic solution

Applications of fixed point theorems to prove the existence of solutions of nonlinear problems usually suffer from the drawback that the uniqueness of solutions has to be proved often by other methods. We will use the following result of Smith [11] which is a slight variation of his Theorem 3.5 (see also Busenberg and Cooke [1]).

THEOREM (Smith). *Let N be an operator on a cone K in the Banach space B and let* $K^* = \{(x_1, x_2) \in K \cap B, 0 \; x1,x2 \in C*$ *where* $C*$*some positive constant is. Assume the following:*

(d) If $(x_1, x_2) \in K^* - \{(0,0)\}$ and $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

| | |
|---|---|
| (e) | N is monotonic nondecreasing on $K^*$. |
| (f) | If $(x_1, x_2)$ is in the interior of $K^*$ and if $0 < \lambda < 1$ then there exists an $\eta > 0$ depending on $(x_1, x_2)$ and $\lambda$ such that N $(\lambda x_1, \lambda x_2) \geq \lambda (1 + \eta)$ N$(x_1, x_2)$. Then N has at most one fixed point different from $(0,0)$ in $K^* \subset B$. |

The cone K= $K_1 \cap K_1$ where $K_1 = \{h \in B_\omega^\tau : h(t) \geq 0, t \in [-\tau, \infty)\}$ is slightly too wide in that solutions of the form $\{x_1(t), 0\}$ and $\{0, x_2(t)\}$ for (2.1) are also possible. We choose a narrower cone $K^*$ in $B$ as follows: define

$$S = \{(x_1, x_2) \in B : b_{1l}/a_{1l} \geq \|x_1\| \geq b_1^u/a_1^u,\ b_{2l}/a_{2l} \geq \|x_2\| \geq b_2^u/a_2^u,\}, \quad \textbf{(3.1)}$$

And let

$$K^* = \{\lambda(x_1, x_2) \in B : (x_1, x_2) \in S, \lambda \geq 0\}. \quad \textbf{(3.2)}$$

Since $\hat{G}_2 / G_1 \subset K^*$, it is not difficult to see that N: $\hat{G}_2 / G_1 \to B$ has at least one fixed point which is a periodic solution of (2.1)-(2.4). We will prove that this solution is unique by applying the result of Smith cited above. The following lemmas are needed first.

LEMMA 4. If $(x_1, x_2) \in K^* - \{0.0\}$ and if $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

PROFF. If $(x_1, x_2) = N(x_1, x_2)$ and $(x_1, x_2) \in K^* - \{0,0\}$, then $(x_1, x_2)$ is a periodic solution of (2.1) –(2.4) and hence we have

$$\begin{aligned} (b_{1l} - a_{1l}x_1)x_1 &< \frac{dx_1}{dt} < (b_1^u - a_1^u x_1)x_1, \\ (b_{2l} - a_{2l}x_2)x_2 &< \frac{dx_2}{dt} < (b_2^u - a_2^u x_2)x_2, \end{aligned} \quad (\mathbf{3.3})$$

And hence any nonzero periodic solution $(x_1, x_2)$ of (2.1)-(2.4) has to be such that

$$b_{1l}/a_{1l} < x_1(t) < b_1^u - a_1^u,$$

$$b_{2l}/a_{2l} < x_2(t) < b_2^u - a_2^u,$$

For $t > -\tau$, which proves the lemma.

LEMMA 5. The mapping $N$ defined by (2.5)-(2.9) is monotone nondecreasing on $K^*$.

PROOF. Let $(x_1, x_2) \in K^*$, $(y_1, y_2) \in K^*$ such that $(x_1 - y_1, x_2 - y_2) \in K^*$ and $(x_1, x_2) \neq (y_1, y_2)$. Then $x_1 \geq y_1$, $x_2 \geq y_2$. From the definition of $N$ we derive that

$N_1$ (t, $x_1, x_2$) – $N_1$ (t, $y_1, y_2$)

$$= \int_0^{\omega} G_{11}(m, \omega, t-s)\{m[x_1(s)] + x_1(s) f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s) f_1(s, y_1(s), y_2(s-\tau_{12}))\} ds$$

$$+ \int_0^t G_{12}(m, 0, t-s)$$

$$\times \{m[x_1(S) - y_1(s)] + [x_1(s) f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s) f_1(s, y_1(s), y_2(s-\tau_{12}))]\} ds \quad \textbf{(3.4)}$$

With a similar relation for $N_2(t, x_1, x_2) - N_2(t, y_1, y_2)$. Since $f_1$ is continuous, $f$ is bounded on $K^*$; also it will follow from $(x_1, x_2) \neq (y_1, y_2)$ that

$$0 \not\equiv x_1(s) - y_1(s) \geq 0 \text{ and } 0 \not\equiv x_2(s) - y_2(s) \geq 0;$$

A sufficiently large positive constant $m$ will make the right side of (3.4) positive.A similar argument will lead to $N_2(t, x_1, x_2) - N_2(t, y_1, y_2) > 0$ perhaps with a larger constant $m$. Finally we can choose the bigger of these two constants $m$ to yield

$$N(x_1, x_2) - N(y_1, y_2) \geq 0 \text{ for } (x_1 - y_1, x_2 - y_2) \in K^* - \{(0,0)\}.$$

The monotonicity of $N$ on $K^*$ follows.

LEMMA 6. Let $(x_1, x_2) \in$ interior of $K^*$ and let $\lambda$ be a constant such that $0 < \lambda < 1$. Suppose $f_1, f_2$ satisfy (2.2) - (2.4). Then there exists a Positive constant $\eta$ such that

$$N(\lambda x_1, \lambda x_2) \geq \lambda(1+\eta) N(x_1, x_2). \quad \textbf{(3.5)}$$

PROOF. A direct computation leads to

$$N_1(t,\ \lambda x_1,\ \lambda x_2) - \lambda N_1(t, x_1,\ x_2)$$

$$= \lambda \int_0^{\omega} G_{11}\ (m, \omega, t-s) g_1\big(s, x_1(s),\ x_2(s-\tau_{12})\big) H_1(s, \lambda,\ x_1,\ x_2) ds$$

$$+ \lambda \int_0^{t} G_{11}\ (m, 0, t-s) g_1\big(s, x_1(s),\ x_2(s-\tau_{12})\big) H_1(s, \lambda,\ x_1,\ x_2) ds$$

**(3.6)**

Where

$$H_1(s, \lambda,\ x_1,\ x_2) = \frac{m\ x_1(s) + x_1(s) f_1\big(s, \lambda\ x_1(s), \lambda\ x_2(s-\tau_{12})\big)}{m\ x_1(s) + x_1(s) f_1\big(s, x_1(s), x_2(s-\tau_{12})\big)} - 1. \qquad (\mathbf{3.7})$$

Since $f_1$ is decreasing (see (2.2)) with respect to $x_1$ and $x_2$ there exists a positive constant say $\eta_1$ such that

$H_1(s, \lambda,\ x_1,\ x_2) \geq \eta_1$ for $0 < \lambda < 1, s > 0, (x_1,\ x_2) \in K^*$, **(3.8)**

And hence

$$N_1(t, \lambda x_1,\ \lambda x_2) \geq \lambda(1+\eta_1) N_1(t, x_1,\ x_2).$$

A similar argument will show that there exists an $\eta_2 > 0$ such that

$$N_2(t, \lambda x_1,\ \lambda x_2) \geq \lambda(1+\eta_2) N_2(t, x_1,\ x_2),$$

And the result follows if we choose $\eta = \min(\eta_1, \eta_2)$.

Now the following uniqueness result is a direct consequence of the result of Smith [11] and the Lemmas 4, 5, 6.

THEOREM 2. *Suppose $f_1, f_2$ of (2.1) satisfy (2.2) – (2.4). Then there exists a unique $\omega$ - periodic strictly positive periodic solution of (2.1).*

## 4. Linear stability of the periodic solution

Let $\{x_1^*(t), x_2^*(t), t \geq -\tau\}$ be the strictly positive unique periodic solution of (2.1). Let $\{x_1(t), x_2(t), t \geq -\tau\}$ be any other solution of (2.1). To study the linear stability of $(x_1^*, x_2^*)$ we let

$$x_1(t) \equiv x_1^*(t)\{1 + X_1(t)\}, \quad x_2(t) \equiv x_2^*(t)\{1 + X_2(t)\}, \quad t \geq -\tau, \tag{4.1}$$

And derive that the perturbations $(x_1, x_2)$ are governed by the linear variational system (after neglecting the non linear terms in $x_1, x_2$)

$$\frac{dX_1(t)}{dt} = -a_{11}(t)\, x_1^*(t) X_1(t) - a_{12}(t) x_2^*(\mathrm{t} - \tau_{12})\, X_2(\mathrm{t} - \tau_{12})$$
$$\frac{dX_2(t)}{dt} = -a_{21}(t) x_1^*(t - \tau_{21}) x_1(t - \tau_{21}) - a_{22}(t) x_2^*(t) x_2(t), \tag{4.2}$$

For $t > 0$, where

$$\left.\begin{aligned} -a_{11}(t) &= \frac{\partial f_1}{\partial x_1} \\ -a_{12}(t) &= \frac{\partial f_1}{\partial x_2} \end{aligned}\right\} \text{Evaluated at } (t, x_1^*(t), x_2^*(\mathrm{t} - \tau_{12})), \tag{4.3}$$

$$\left.\begin{aligned} -a_{22}(t) &= \frac{\partial f_2}{\partial x_1} \\ -a_{22}(t) &= \frac{\partial f_2}{\partial x_2} \end{aligned}\right\} \text{Evaluated at } (t, x_1^*(t - \tau_{21}), x_2^*(\mathrm{t})).$$

To establish the linear asymptotic stability of $(x_1^*(t), x_2^*(\mathrm{t}))$ it is enough to prove that of the trivial solution of (4.2) - (4.3) We will do this by means of a combination of a comparison differential system and a Lyapunov function. First we define

$$\frac{d|X_1(t)|}{dt} = \begin{cases} \frac{dX_1(t)}{dt} & if\ X_i(t) \geq 0, \\ -\frac{d\backslash X_1(t)}{dt} & if\ X_i(t) < 0, \end{cases} \quad i = 1,2, \tag{4.4}$$

And derive from (4.2)- (4.3) the following comparison system of differential inequalities:

$$\frac{d|X_1(t)|}{dt} \leq -b_{11}x_1^*(t)|X_1(t)| + b_{12}x_2^*(t-\tau_{12})|X_2(t-\tau_{12})|,$$
$$\frac{d|X_2(t)|}{dt} \leq -b_{21}x_1^*(t-\tau_{21})|X_1(t-\tau_{12})| - b_{22}x_2^*(t)|X_2(t)|, \quad \textbf{(4.5)}$$

For $> 0$, where

$$-b_{11} = \max_{t\in-\tau}\{-a_{11}(t)\},$$

$$b_{12} = |\min_{t\in-\tau}\{-a_{12}(t)\}|$$
$$b_{21} = |\min_{t\in-\tau}\{-a_{21}(t)\}| \quad \textbf{(4.6)}$$

$$-b_{22} = \max_{t\in-\tau}\{-a_{22}(t)\}$$

We can now formulate the result as follows. By the periodicity of $a_{ij}$, the constants $b_{ij}$ are independent of $\tau$ (see also (2.2)).

THEOREM 3. *The unique strictly positive periodic solution* $\{x_1^*(t), x_2^*(t), t\in-\tau\}$ of $(2.1)-(2.4)$ *is linearly asymptotically stable provided*

$$|-b_{11}| > b_{21} \text{ and } |-b_{22}| > b_{12}, \quad \textbf{(4.7)}$$

*Where the constants* $b_{ij}(i,j=1,2)$ *are as defined in* (4.6).

PROOF. Consider the scalar valued Lyapunov-like function $v(t) = V(t, X_{1,}X_{2,})$where

$$v(t) = V(t, X_{1,}X_{2,}) = |X_{1,}(t)| + |X_{2,}(t)| + \int_{t-\tau_{12}}^{t} b_{12}x_2^*(u)|X_{2,}(u)|du$$
$$+\int_{t-\tau_{21}}^{t} b_{21}x_1^*(u)|X_{1,}(u)|du. \quad \textbf{(4.8)}$$

Since $x_1^*, x_2^*$ are strictly positive, it is found from (4.7)-(4.) that

$\frac{dv}{dt} \leq -[(|b_{11}| - b_{21})\, x_1^*(t)|X_{1,}(t)| + (|b_{22}| - b_{22})\, x_2^*(t)|X_2(t)|]$ **(4.9)**

For $t > 0$. It follows from (4.7)-(4.9) that $V(t, X_{1,}X_{2,}) \leq |X_{1,}(t)| + X_{2,}(t)|$ and $dv/dt \leq 0$,which will imply the stability of the trivial solution of (4.2)-(4.3). We claim that (4.9) also implies the asymptotic stability of the trivial solution of (4.2)-(4.3)in the sense that

$$\lim_{t-\infty} [\,|X_{1,}(t)| + |X_{2,}(t)|\,] = 0,$$

**(4.10)**

For all bounded initial conditions $X_{1,}(s), X_{2,}(s), s \in \{\mathcal{T}, 0\}. Suppose (4.10) is\ not$ true Then we can find an increasing sequence $\{t_m\} \to \infty$ as $m \to \infty$ such that

$(\,|b_{11}| - b_{21})\, x_1^*(t_m)|X_{1,}(t_m)| + (|b_{22}| - b_{12} x_2^*(t_m)|X_2(t_m)| > \varepsilon,$

M=1, 2, 3…

**(4.11)**

Using the right derivative *dv/dt* of $v$ we derive that

$$v(t_m - \varepsilon) - v(t_m) \leq -\int_{t_m-\varepsilon}^{t_m} \frac{dv(T)}{dT}\, d\tau.$$

**(4.12)**

If $\varepsilon$ is small enough we have from (4.9)(4.11),(4.12)that $dv(\tau)/dr < (-\varepsilon/2)$ and

$v(t_m - \varepsilon) - v(t_m) \geq \varepsilon^2/2$ and hence

$v(t_m) \geq v(t_m - \varepsilon) - (\varepsilon^2/2)$

$\geq v(t_m - 1)$-$(\varepsilon^2/2)$

$\geq v(t_m - 2)$-$(\varepsilon^2/2)$

$v(t_0)$-m ( $\varepsilon^2/2$)

Thus for sufficiently large $m, v(t_m)$becomes negative which contradicts the Nonnegative nature of $v\ for\ t$ 0. this contradiction shows that $|X_{1,}(\mathrm{t})| + |X_2(\mathrm{t})| \to 0$ as $\mathrm{t} \to \infty$ and the proof is complete.

As an example let us consider the case of two species Lotka- Volterra type competition in a periodic environment for for which we have

$$f_1\big(\mathrm{t}, x_1(\mathrm{t}), x_2(\mathrm{t}-T_{12})\big)1 \quad \left\{ \frac{x_1(\mathrm{t}) + ax_2(\mathrm{t}-T_{12})}{K_1(\mathrm{t})} \right\}$$

$$f_2(\mathrm{t}, x_1(\mathrm{t}-\ {}_{21})x_2(\mathrm{t})) = 1 - \left\{ \frac{\beta x_1(\mathrm{t}-T_{21})+x_2(\mathrm{t})}{K_2(\mathrm{t})} \right\}$$

Where $\alpha, \beta$ are positive constants and $k_{1,}k_{2,}$ are strictly positive continuous periodic functions with a common period Equation (4.6)and (4.7) for this special case lead to

$$\left|\mathrm{axm}_{t>T}\{-a_{11}(t)\}\right| = \max_{t>T}(1/K_1(t)) > \min_{t>T}\big(-a_{21}(\mathrm{t})\big)\Big| = \min_{t}, \beta/K_2(\mathrm{t})$$

$$\max_{t T}\{-a_{22}(t)\}\Big| = \max_{t T}(1/K_2(t)) > \min_{t T}\big(-a_{12}(\mathrm{t})\big)\Big| = \min_{t} \ \alpha/K_1(\mathrm{t})$$

The inequalities reduce in the absence of temporal variations in $k_{1,}k_{2,}$ to the inequality 0< $\alpha,\beta$<1, which is well known for the global asymptotic stability of the positive steady state corresponding to the above example in a constant systemic domain the above example shows also that delays in interspecific interactions have no influence on the local asymptotic stability of the unique periodic solution of the periodic ecosystem A number of cases of this type in a constant environment have been established recently by the author (Gopalsamy [6]).A numerical example is illustrated at the end.

**Consequences**

It can be found that all the conclusions on the existence uniqueness and stability of a periodic solution of (2.1)-(2.4) are valid when $OT_{12}=0$ or $T_{12}=0$ or both $T_{12}=0$ and $T_{21}=0$: thus it will follow that delays in interspecific competition do not destabilize an otherwise stable coexistence (oscillatory or nonoscillatory) of two competing types of interacting particles atoms molecules transactions interaction. In forthcoming work we will extend the analysis of the system (2.1)-(2.4) to a system of n-competing species with a hypothesis similar to that of (2.4) It is however an open problem to investigate the effects of delays in intraspecific competition in a periodic milieu

Pianka ([8],pages 139-140) remarks that "In real population rate of increase competitive abilities and carrying capacities do very from individual to individual with population density in space and time Indeed temporal variation in the environment may often allow coexistence by continually altering the competitive abilities of population inhabiting it Moreover time lags are doubt lessly of some importance in real population Finally a heterogeneous environment may allow real competitors to evolve divergent resource utilization patterns and to reduce interspecific competitive inhibition

To understand the mechanism of selection pressure against competitive inhibition we first write the Malthusian parameters $\mu_1, \mu_2$ of fitness for the two species as follows (see MacArthur[7],page 226-230):

$$\mu_1=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_1\left(t,x_1(t),x_2(t-T_{12})\right)$$

$$\mu_2=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_2(t,x_1(t-T_{21}),x_2(t))$$

Since $\left(x_1(t),x_2(t)\right)\rightarrow\left(x_1^*(t),x_2^*(t)\right)$ for large t we can write the parameters of finess in the from

$$\mu_1=f_1\left(t,x_1^*(t),x_2^*(t-T_{12})\right)$$

$$\mu_2=f_2\left(t,x_1^*(t-T_{12}),x_2^*(t)\right)$$

In these relations we can use

$$x_2^*(t-T_{12}), x_2^*(t) - T_{12}\frac{dx_2^*(t-T_{12}+\theta_1 T_{12}}{dt}$$

$$x_1^*(t-T_{12}), x_1^*(t) - T_{21}\frac{dx_1^*(t-T_{21}+\theta_2 T_{21}}{dt}$$

$\theta_1$,$\theta_2$,$\in (0,1)$ and rewrite $\mu_1$,$\mu_2$ in the form

$$\mu_1 = f_1\left(t, x_1^*(T)\right), x_2^*(t) - T_{12}\frac{dx_2^*(t-T_{12}+\theta_1 T_{12}}{dt}$$

$$\mu_2 = f_2\left(t, x_1^*(T)\right), -T_{21}\frac{dx_1^*(t-T_{21}+\theta_2 T_{21}}{dt}, x_2^*(t))$$

$\theta_1$,$\theta_2$,$\in (0,1)$

It is seen that the presence of the delays $T_{12}$ and $T_{21}$ provides additional sources of variability in the finess parameters and thereby in the selection pressure against interspecific competition Since a necessary and sufficient condition for natural selection to act is differential (variable) rates of confederation growth we find that time delays coupled with periodicity of an domain may allow real to evolve divergent resource utilization patterns and to reduce

Interspecific competitive inhibition for some remarks relating to the significance of a periodic environment to sustain the coexistence of a large number of confederation types in the absence of delays we refer to a recent article (Gopalsamy)The following example has been numerically solved and the solutions are displayed graphically for two different values of the delay parameter $\mathcal{T}$:

$\frac{dx_1(t)}{dt} x_1(t)[1 - x_1(t) + 0.4x_2(t-\mathcal{T})/k_1(t)]$

$\frac{dx_2(t)}{dt} x_2(t)[1 - \{0.8\, x_1(t-\mathcal{T}) + x_2(t)\}/k_2(t)]$

Where $k_1(t) = 20 + \sin(\pi t)$ and $k_2(t)$=2.5+sin ($\pi$t)

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There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. “All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ....” (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To “deconstruct” is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between “speech” and “writing.” This binary logic functions in an illicit way to establish speech as the means of giving “presence” to the world, while writing is deemed derivative and inferior In Derrida’s sense of “grammatology,” however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of “differance” (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons

and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti

thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification.

## GOPALSAMY'S MODEL APPIED TO COSMOLOGICAL PROBLEMS

## GLOBAL ASYMPTOTIC STABILITY IN AN ALMOST PERIODIC LOTKA- VOLTERRA SYSTEM

**Abstract**

**Gopalsamy studied** Sufficient conditions are obtained for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of a Lotka- Volterra system with almost periodic coefficients. Research was published Australian. Math Soc. B.27 (1986) 346-360. Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of decay rate of perturbations. Sole intention, main objective and primary concern of the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are counted in the same vein as that of predator prey. Note that the model is applicable to each and every module. Quintessentially each parameter should have been given different name such as superscript 1,2, etc., as that has been done in the foregoing and a formal concatenation done at the end. Due to spatial restrictions we have not been able to do the line of thought and manner of analysis as has been done hereinbefore. This is the most important aspect of the model. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated

model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 3. Introduction

At the outset we assume that the system. This article is concerned with the derivation of a set of "easily verifiable" sufficient conditions for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of the Lotka- Volterra system

$$\frac{dx_1(t)}{dt} = x_i(t)\left\{b_1(t) - \sum_{j=1}^{n} a_{ij}\,(t)x_j(t)\right\};$$

$$t>t_0;\ t_0 \in (-\infty,\infty);\ i = 1,2,\dots,n, \qquad (1.1)$$

where $b_1, a_{ij}\ (i,j = 1,2,\dots.,n)$ are nonnegative almost periodic functions defined for $t \in (-\infty,\infty)$.

In question satisfies Haimovici's equation based on age wise stratification given in the foregoing. Note that Haimovici's equations can be applied to Lotka-Volterra system term by term.

Let R and $R_n$ respectively denote the set of all real numbers and the $n$ - dimensional real Euclidean space; $R_n^+$ will denote the nonnegative cone of $R_n$. Since almost- periodic functions defined on R are bounded, we define the constants $b_i^l, b_l^u, a_{ij}^l, a_{ij}^u, (i,j = 1,2,\dots,n)$ by the following:

$$\inf_{t\in R}\ b_l(t) = b_i^l; \qquad \sup_{t\in R} b_l\,(t) = b_l^u, \qquad i,j = 1,2,\dots,n.$$

$$\inf_{t\in R} a_{lj}\,(t) = a_{lj}^l; \qquad \sup_{t\in R} a_{lj}\,(t) = a_{ij}^u.$$

We shall study the almost- periodic system (1.1) with the following assumption on the coefficients of (1.1);

$$b_l^l > 0; \qquad a_{ll}^l > 0, \qquad (1.2)$$

$$b_l^l > \sum_{\substack{j=1 \\ j\neq i}}^{n} a_{lj}^u\,(b_j^u\ /\ a_{lj}^l). \qquad (1.3)$$

Since the solutions of (1.1) corresponding to nonnegative initial conditions remain nonnegative subsequently, it will follow that

$$\frac{dx_1}{dt} \geq x_l\{b_l^u - a_{ll}^l x_i\}; \qquad t > t_0, i = 1, 2, \dots, n, \qquad (1.4)$$

As a consequence of which we will have

$$0 < x_l(t_0) \leq b_l^u / a_{ll}^l = x_i^u \Rightarrow x_i(t) \leq b_l^u / a_{ll}^l, \qquad i = 1, 2, \dots, n; t \geq t_0. \qquad (1.5)$$

(1.2) And (1.5) together lead to

$$\frac{dx_1}{dt} \geq x_l \left\{ b_l^u \sum_{\substack{j=1 \\ j \neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right) - a_{ll}^l x_i \right\}, \qquad i = 1, 2, \dots, n ; \qquad (1.6)$$

Now (1.2), (1.3) and (1.6) lead to

$$x_i(t_0 \geq \left( b_i^l - \sum_{\substack{J=1 \\ J \neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right) \right) (1/a_{ll}^u) = x_i^l$$

$$\Rightarrow x_i(t) \geq x_i^l \ for\ t \geq t_0,\ i = 1, 2, \dots, n. \qquad (1.7)$$

4. It is our aim to show that the mode together with other models stated in the foregoing together with haimovicis model is applicable to whole lot of consummative debit – credit systems in cosmology, astrophysics, finance, economics, sociology, anatomy computation. We will not however divide the system studied here into 9 components/ modules due to restrictions on space….

From the foregoing preparation we have the following:

LEMMA 1.1. *Assume that* $b_i, a_{ij} (i, j = 1, 2, \dots, n)$ *are scalar nonnegative almost periodic functions defined on* $\boldsymbol{R}$ *such that (1.2) and (1.3) hold. Then the set S defined by*

$$S = \{x = (x_1, x_2, \dots, x_n) \in \boldsymbol{R}_n | x_* \leq x_l \leq x^*; i = 1, 2, \dots, n\}$$

(1.8)

*Where*

$$x_* = \min_{1} x_l^l ; \qquad x_* = \min_{1} x_i^u \qquad (1.9)$$

*is invariant with respect to (1.1)*

**Appellation, cognomen, compellation, designation, denomination of the conducive, infelicitous and unfortutious** responses to environment change may be influenced by changes in available domain and milieu under consideration. An example is the following study: Authors describe the gravitational degrees of freedom of the Schwarzschild black hole by one free variable. Authors introduce an equation which we suggest to be the Schrödinger equation of the Schwarzschild black hole corresponding to this model. We solve the Schrödinger equation explicitly and obtain the mass spectrum of the black hole as such as it can be observed by an observer very far away and at rest relative to the black hole. Our equation implies that there is no singularity inside the Schwarzschild black hole, and that the black hole has a certain ground state in which its mass is non-zero. **Schrodinger equation of the Schwarzschild black hole Jarmo Mäkelä (Cambridge U E-Print: gr-qc/9602008 | PDF** Particles, axions, atoms, molecules, quarks, Majorana fermions, bosons et al in its confederation capacity and aggregation processes, type interactions and interactions between landscape dynamics. Following bears ample testimony, infallible observatory to the fact there might be changes in the quantum behaviour due to merger of black holes. Following example bears enough testimony, infallible observatory and impeccable demonstration and apodictic evidence as how some of the models could be applied to cosmological problems. This is attributable and ascribable to the fact that the classification criterion is based on the age.

**How gravity kills (e) Schrödinger's cat Theorists argue that warped space-time prevents (e) quantum superpositions of large-scale objects. Elizabeth Gibney 17 June 2015 Propositions are written in model able form and can easy be scheduled.**

(23) If the cat in Erwin Schrödinger's famous thought-experiment behaved according to quantum theory, it would be able to exist in multiple states at once: both dead and alive.

(24) Physicists' common explanation for why we don't (e) see such quantum superpositions — in cats or any other aspect of the everyday world — is interference from (e) the environment.

(25) As soon as a quantum object interacts with (e&eb) a stray particle or a passing field, it picks (e) just one state, collapsing (eb) into our classical, everyday view.

(26) But even if physicists could completely isolate (eb) a large object in (eb) a quantum superposition, according to researchers at the University of Vienna, it would still collapse into (e&eb) one state — on Earth's surface, at least.

(27) "Somewhere in interstellar space it could be that the cat has (e) a chance to preserve quantum coherence, but on Earth, or near any planet, there's (=) little hope of that," says Igor Pikovski. The reason, he asserts, is gravity.

(28) Pikovski and his colleagues' idea, laid out in a paper published in Nature Physics on 15 June1, is at present only a mathematical argument. But experimenters hope to test whether gravity really does collapse (eb) quantum superpositions, says Hendrik Ulbricht, an experimental physicist at the University of Southampton, UK. "This is a cool, new idea, and I'm up for trying to see it in experiments," he says. Assembling the technology to do so, however, may take as long as a decade, he says.

(29) Cinema-goers who saw the film Interstellar are already familiar with the basic principle behind the Vienna team's work. Einstein's theory of general relativity states that an extremely massive object causes clocks near it to (e) run more slowly because (e) its strong gravitational field stretches (e&eb) the fabric of space-

time (which is why a character in the film aged only an hour near a black hole, while (e) seven years passed on Earth). On a subtler scale, a molecule placed nearer the Earth's surface experiences (e&eb) a slightly slower clock than one placed slightly further away.

(30) Because of gravity's effect on (e&eb) space-time, Pikovski's team realised that variance in a molecule's position will also influence (e&eb) its internal energy — the vibrations of particles within (eb) the molecule, which evolve (eb) over time.

(31) If a molecule were put in a quantum superposition of two places, the correlation between position and (e&eb) internal energy would soon cause (eb) the duality to 'decohere' to (e) the molecule taking just one path, they suggest.

(32) "In most situations decoherence is due to (e) something external; here it's as though the internal jiggling is interacting with (e&eb) the motion of the molecule itself," adds Pikovski.

**A practical limit**

(33) No one has yet seen this effect because other sources of decoherence — such as magnetic fields, thermal radiation and vibrations — are (=) typically much stronger and cause (eb) quantum systems to collapse long before gravity becomes an issue. But experimenters are keen to try.

(34) Markus Arndt, an experimental physicist also at the University of Vienna, already tests whether quantum superpositions can be observed for (e) large, though not cat-sized, objects (see Fattening up Schrödinger's cat).

(35) He sends large molecules through (e&eb) a matter-wave interferometer, a system that gives (eb) each molecule a choice of (e) two different paths.

(36) In the classical view, a molecule travels down just one path; a quantum molecule effectively passes through (e&eb) both routes at once and interferes with (e&eb) itself to create (eb) a characteristic wave-like pattern (see image, below).

(37) A quantum interference pattern built up from complex molecules known as phthalocyanines.(deleted due to spatial constraints)

(38) A similar set up could be used to (e) test for gravity's ability to destroy (e) quantum behaviour: by comparing a vertical interferometer, in which the superposition should soon decohere due to (e) the stretching of time in one path versus (e&eb) the other, with a horizontal set up, where (e) the superposition could remain.

(39) Arndt, who has tested the effect for (e) molecules of up to 810 atoms2, points out those large molecules, would be (=) good for testing the gravitational effect because (e) they contain (e) lots of particles contributing to (e) internal energy.

(40) **But not only would the researchers have to further suppress the external environment to reduce (e) other decoherence effects, they would also need to (e) increase the separation of the two paths from micrometres to metres, or else use molecules perhaps 1 million times more massive. "It's certainly very challenging," says Arndt**.

(41) If gravity's effect does limit (e) quantum behaviour on Earth, tests of quantum reality for (e) large objects may eventually have to (e) move into space, says Angelo Bassi, a physicist at the University of Trieste in Italy.

(42) "But from a deep, fundamental point of view, this is nothing new," he says. A gravitational field is merely another environment to interact with, so invoking it does not (e) explain whether quantum behaviour might lead to (eb) classical reality if gravity's influence were mitigated — for example, by doing the experiment in (eb) gravity-free space.

(43) Nor does the effect described by Pikovski and colleagues say anything about quantum gravity: a theory that would unite gravity and (e&eb) quantum mechanics in (eb) a single description, which many researchers are working towards.

(44) "It's an interesting effect, but it's still quantum physics applied to (e&eb) classical general relativity. In that way it doesn't change our picture of the world," adds Bassi. Nature doi:10.1038/nature.2015.17773

Integrated factors these influences because they deal with ambience, mise en scene change or confederation dynamics, but rarely both. Time series of (e) space time suitability models with spatially explicit stochastic models to explore factors that influence the viability of type viability and confederation feasibility under stable and changing scenarios. Methodology could be applied to interrogation propositions and problems which might subsume one with the other notwithstanding that there is a space of singularities and there is space for completion of present from the past and future. An example is Time image movement from Deleuze. We clearly face a stiff challenge in our attempt to proffer Deleuzian treatments of photography as our critical project brushes up against the fact that Deleuze does not invest photography with (e&eb)the creative and re-creative potential that he associates with (e&eb) both cinema and modern painting. In short, while he highly values other visual artistic forms, he seemingly presents photographic texts as stagnate documents or tools that produce certainty, organize bodies and desires, and iterate hackneyed ideas. He even uses photography as something of a foil to demonstrate the innovation of the cinema and the originality of modern painters. In Francis Bacon: The Logic of Sensation (1981), he notes how "photography has taken over the illustrative and documentary role, so that modern painting no longer needs to fulfill this function, which still burdened earlier painters" (10). He later explains that "photographs are ways of seeing, and as such, they are illustrative and narrative reproductions or representations ... . They are what are seen, until finally one sees nothing else." He treats the photo as an instrument for reproducing representations of reality—a device that iterates images until they are ossified as established stories, icons, or even stagnant perceptions. Indeed, like Bergson before him, Deleuze compares photography to what he identifies as a reductionist model of perception: "what we see, what we perceive, are photographs" (74).[4] The photo is, in effect, an always-already passé sensational experience that numbs our sensitivity to the ongoing vitality of life and provides us with an efficient understanding of the dynamism of our world. But it is compelling that Deleuze theorizes both the quotidian nature and the liberatory potential of photography; he treats it as a common, everyday visual experience that, despite its aesthetic and technological limitations, emancipates the modern painter to explore new artistic opportunities. This special issue of Rhizomes repeatedly provides artistic examples and critical strategies that invite us to see photography as replete with the power to unburden and emancipate that Deleuze associates with both film and painting. Indeed, one way to think of this project of reading Deleuze against Deleuze might be to ask: if photography can free painting and painters to pursue novel creative opportunities, can it likewise free itself, photographic practitioners, and even viewers to create and re-create anew? **Rhizomes » Issue 23 (2012) » Michael Kramp Unburdening Life, or the Deleuzian Potential of Photography Michael Kramp Lehigh University** dmk209@lehigh.edu

**Existence of an almost- periodic solution**

DEFINITION 2.1. A function $f(t,x)$,where $f$ is an $m$- vector, $t$ is a real scalar and $x$ is an $n$- vector, is said to be almost periodic in $t$ uniformly with respect to $x \in X \subset R_n$, $if\ f(t,x)$ is continuous in $t \in R\ and\ x \in X$, and if for any $\varepsilon > 0$, it is possible to find a constant $l(\varepsilon) > 0$ such that in any interval of length $l(\varepsilon)$ there exists a $\tau$ such that the inequality

$$\left||f\ (t+\ \tau,x)-\ f(t,x)|\right|=\sum_{l=1}^{m}|f_l\ (t+\ \tau,x)=\ f_l(t,x)|$$

$$<\varepsilon \qquad (2.1)$$

is satisfied for all $t\in(-\infty,\infty), x\in X$. The number $\tau$ is called an $\varepsilon-$ translation number of $f(t,x)$.

Let $\{\lambda_j\}$ denote the set of all real numbers such that

$$\lim_{T\to\infty}T^{-1}\int_0^T f(t,x)\exp(-i\,\lambda_j t)dt\ \not\equiv 0\quad for\ x\ \in X; i=\sqrt{-1}. \qquad (2.2)$$

It is known that when the set X is separable, the set of numbers $\{\lambda_j\}$ in (2.2) is countable. The set $\{\sum_1^N n_j\ \lambda_j\}$ for all integers N and integers $n_j$ is called the module of $f(t,x)$.

The following result concerned with module containment will be used in the proof of our existence theorem below.

**MODULE CONTAINMENT THEOREM** (Yoshizawa [7], page 18). *Let $f(t,x)$ and $g(t,x)$ be almost periodic in t uniformly for $x\ \in D\subset\ \boldsymbol{R}_n$. If for any compact set $S\subset D$ and for any sequence of real numbers $\{\tau_k\}$ having a limit as $k\to\infty$ (the limit being finite or$\pm\infty$) for which $\{f(t+\tau_k,x)\}$ is uniformly convergent on $\boldsymbol{R}\times S, \{g(t+,x)\}$ is also uniformly convergent on $\boldsymbol{R}\times S$ then the module of $g(t+,x)$ is contained in the module of $f(t+,x)$.*

We can now formulate our existence theorem.

**THEORM 2.1** *Assume that the almost-periodic parameters $b_1, a_{1J}\ (i,j=1,2,\dots,n)$ of (1.1) satisfy (1.2)-(1.3). Furthermore assume that there exists a positive number $\mu$ such that*

$$\inf_{1\in\boldsymbol{R}}a_{11}(t)>\sum_{\substack{J=1\\J\neq 1}}^{n}(\sup_{t\in\boldsymbol{R}}a_{Jl}(t))+\mu,\quad i$$

$$=\ 1,2,\dots,n. \qquad (2.3)$$

*Then the system (1.1) has a strictly positive (component wise) almost-periodic solution, say $w\ (t)=w_1(t),\dots,w_n(t), t\in\boldsymbol{R}$, whose module is contained in that of*

$$f(t,x)=\{f_1(t,x_1,\dots,x_n),f_2(t,x_1,\dots,x_n),\dots,f_n(t,x_1,\dots,x_n)\}$$

*Where*

$$f_l(t, x_1, \dots, x_n) = x_l \left\{ b_l(t) - \sum_{j=1}^{n} a_{lj}(t) x_j \right\},$$

$$i = 1, 2, \dots n; (x_1, x_2, \dots, x_n) \in S.$$

(2.4)

PROOF. The set $S$ defined by (1.8)-(1.9) is compact in $\boldsymbol{R}_n$. Let $\{\tau_h\}$ be any sequence of real numbers such that $\tau_k \to \infty$ as $k \to \infty$ and

$$\sum_{l=1}^{n} |f_i(t + \tau_k, x_1, \dots, x_n) - f_l(t, x_1, \dots, x_n)| \to 0 \quad \text{as } k \to \infty \qquad (2.5)$$

Uniformly in $(x_1, \dots, x_n) \in S$ for $t \in (-\infty, \infty)$. for any real number $\beta$let $k_0(\beta)$ be the smallest value of $k$ such that $\tau k_0 + \beta \geq t_0$. Since the invariance of $S$ implies that any solution of (1.1)

$$\{x_1(t + \tau_k, t_0, x_1(t_0)), x_2(t + \tau_k, t_0, x_2(t_0)), \dots, x_n(t + \tau_k, t_0, x_n(t_0))\} \in S$$

For all $t \geq \beta$ and $k \geq k_0(\beta)$ whenever $\{x_1(t_0), \dots, n_n(t_0)\} \in S$, our first task in the proof is to show that the sequence

$$\{x_1(t + \tau_k, t_0, x_1(t_0)), \dots, x_n(t + \tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Converges to a continuous bounded function say $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined for $t \in [\beta, \infty)$ such that $w_1(t) \in [x_*, x^*], t \in [\beta, \infty), i = 1, 2, \dots, n$ and the convergence is uniform on all compact subsets of $[\beta, \infty)$. Since $x_*$ and $x^*$ are independence of $\beta$, it is enough to show that

$$\{x_1(t + \tau_k, t_0, x_1(t_0)), \dots, x_n(t + \tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Is a Cauchy sequence on compact subsets of $[\beta, \infty)$.

Let $U$ be any compact subset of $\beta, \infty)$., and let be any arbitrary positive number. Choose an integer $n_0(\varepsilon, \beta) \geq k_0(\beta)$ so large that for $m \geq k \geq n_0(\varepsilon, \beta)$

and all $t \in (-\infty, \infty)$

$$\sum_{i=1}^{n} |b_1(t + \tau_m) - b_1(t + \tau_k)| < (\mu \varepsilon x_*)/(4x^*), \qquad (2.6)$$

$$x^* \sum_{i=1}^{n} \sum_{j=1}^{n} |a_{lj}(t+\tau_m) - a_{lj}(t+\tau_k)| < (\mu\varepsilon x_*)/(4x^*) \quad (2.7)$$

$$(2nx^*/x_*)\exp[-\mu x_*(\beta + \tau_k - t_0)] < (\varepsilon/2x^*)$$

(2.8)

Denote $x_i(t, t_0, x_1, (t_0))$ $by$ $x_1(t)$ brevity, and consider a function $v(s)$ defined by

$$v(s) = v\left(x(s), x(s+\tau_m+\tau_k)\right)$$

$$= \sum_{i=1}^{n} |\log x_1,(s) - \log x_1(s+\tau_m+\tau_k)|, \qquad x \in t_0, \quad (2.9)$$

Where $x(t) = \{x_1(t), \dots, x_n(t)\}$ denotes a solution of (1.1) with $\{x_1,(t_0), \dots, x_n(t)\} \in S$. the invariance of S for (1.1), together with the elementary mean value theorem of differential calculus, implies

$$(1/x^*) \sum_{l=1}^{n} |x_l(s) - x_1(s+\tau_m+\tau_k)| \in v(s)$$

$$\in (1/x^*) \sum_{i=1}^{n} |x_l(s) - x_1(s+\tau_m+\tau_k)|. \quad (2.10)$$

Calculating the right derivative $D^+v$ of $v$ along the solutions of (1.1) and simplifying

$$D^+v(s) = \sum_{i=1}^{n} \{sign[x_i(s) - x_i(s+\tau_m-\tau_k)]\}$$

$$\left\{\frac{d[log x_i(s)]}{ds} - \frac{d[log x_i(s+\tau_m-\tau_k)]}{ds}\right\}$$

$$\in \sum_{l=1}^{n} |b_l(s) - b_l(s+\tau_m-\tau_k)|$$

$$+\sum_{l=1}^{n}(-a_{il}(s)|x_i(s)-x_i(s+\tau_m-\tau_k)|+\sum_{\substack{j=1\\ j\neq l}}^{n}a_{ij}(s)|x_j(s)-x_j(s+\tau_m-\tau_k)|)$$

$$+\sum_{i=1}^{n}\sum_{j=1}^{n}|a_{ij}(s)|a_{ij}(s+\tau_m-\tau_k)|\,x_j(s+\tau_m-\tau_k) \qquad (2.11)$$

$$\leq -\mu x_* v\big(x(s),x(s+\tau_m-\tau_k)\big)+\sum_{l=1}^{n}|b_l(s)-b_i(s+\tau_m-\tau_k)|$$

$$+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|. \qquad (2.12)$$

Let us set $s-\tau_k=\sigma$; then we have from (2.6)-(2.7) that

$$\sum_{i=1}^{n}|b_l(s)-b_1(s+\tau_m-\tau_k)|+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|.$$

$$=\sum_{l=1}^{n}|b_l(\sigma+\tau_k)-b_i(\sigma+\tau_m)|$$

$$+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(\sigma+\tau_k)-a_{ij}(\sigma+\tau_m)|$$

$$<(\mu\varepsilon x^*)/(2x^*). \qquad (2.13)$$

Now an integration of (2.12) over $(t_0,t-\tau_k)$ together with (2.13) leads to

$$v\big(x(t+\tau_k),x(t+\tau_m)\big)\leq v\big(x(t_0),x(t_0+\tau_m-\tau_k)\big)\exp[-\mu x_*(t+t_k-t_0)]$$

$$+[(\mu\varepsilon x_*)/(2x^*)]/\mu x_*)$$

$$\leq[(2x_n^*)/x_*]\exp[-\mu x_*(t+t_k-t_0)]+(\varepsilon/2x^*)$$

$< (\varepsilon/x^*)$ On using (2.8), (2.10) (2.14)

For $m \geq n_0(\varepsilon, \beta)$ and for all $t \in \cup \subset [\beta, \infty]$, Thus we have from the definition of $\upsilon$ and (2.14) that

$$\sum_{l=1}^{n} |log x_i\,(t+\tau_k) - log x_l(t+\tau_m)| < (\varepsilon/x^*) \qquad (2.15)$$

And hence using $x_* \leq x_i(t) \leq x^*, i = 1, 2, \dots, n; t \geq t_0$,

$$\sum_{i=1}^{n} |x_i\,(t+\tau_k) - x_i(t+\tau_m)| < \varepsilon \text{ for all } t \in \cup \subset [\beta, \infty], m \geq n_0(\varepsilon, \beta). \qquad (2.16)$$

The existence of a function $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined on $[\beta, \infty]$, such that $w_i(t) \in [x_*, x^*], i = 1, 2, \dots, n; t \in [\beta, \infty]$ follows from (2.16) and the fact that $x_i(t) \in [x_*, x^*], i = 1, 2, \dots, n;$ for $t \geq t_0$. Since $\beta$ is arbitrary, $w$ is in fact

Defined on $(-\infty, \infty)$ and hence we have

$x_i(t + \tau_k) \to w_l(t)$ as $k \to \infty$ uniformly on compact subsets of $(-\infty, \infty)$, $i = 1, 2, \dots, n.$

(2.17)

Our next task is to show that $w: R \to S$ is differentiable and is a solution of (1.1). Since $x(t) = \{x_1(t), \dots, x_n(t)\}$ satisfies (1.1) we have $\sum_{i=1}^{n} |\frac{d}{dt}\, x_i(t+\tau_k) - \frac{d}{dt}\, x_i(t+\tau_m)|$

$$\leq \sum_{i=1}^{n} |f_l\,(t+\tau_k, x_1(t+\tau_k), \dots, x_n\,(t+\tau_k))$$

$$- f_l(t+\tau_m, x_1(t+\tau_k), \dots, x_n(t+\tau_k))| \qquad (2.18)$$

$$+ \sum_{I=1}^{n} |f_l\,(t+\tau_m, x_1(t+\tau_k), \dots, x_n(t+\tau_k))$$

$$-f_l\big(t+\tau_m, x_1(t+\tau_m), \dots, x_n(t+\tau_m)\big)|$$

Where $f_l\ (i=1,2,\dots,n)$are as in (2.4). Since x$(t+\tau_k) = x_1(t+\tau_k), \dots, x_n(t+\tau_k) \in S$ for large $\tau_k$ for $t$ in each compact subset of $(-\infty,\infty)$ there exists an integer $n_1(\varepsilon) > 0$ such that if $m$ ☒ $k$ ☒$n_1(\varepsilon)$, then for any $\varepsilon > 0$,

$$\sum_{i=1}^{n} |f_l\left(t+\tau_k, x_1(t+\tau_k), \dots, x_n(t+\tau_k)\right)$$

$$-f_1\big(t+\tau_m, x_1(t+\tau_k), \dots, x_n(t+\tau_k)\big)| < (\varepsilon/2). \qquad (2.19)$$

Similarly since $x_l(t+\tau_k) \in S$ and $x_l(t+\tau_m) \in S$ for $i=1,2,\dots,n$ there is an integer $n_2(\varepsilon) > 0$ such that for large enough $m, k$, such that $m$ ☒ $k$ ☒$n_2(\varepsilon)$, we have

$$\sum_{l=1}^{n} |f_l\left(t+\tau_m, x_1(t+\ \tau_k), \dots, x_n(t+\ \tau_k)\right)$$

$$-f_1\big(t+\tau_m, x(t+\tau_m), \dots, x_n(t+\tau_m)\big)| < (\varepsilon/2) \qquad (2.20)$$

This equation will be very much useful in chemistry if used with the modules mentioned and written in concatenation with each other. It should be borne in mind that the applicability is modulewise and in no case concatenation should be marginalized. In all cases parameters are different which could be represented with different superscripts and subscripts like 1,23,…….etc., **Anal Bioanal Chem. 2004 Sep; 380(1):36-45. Epub 2004 Jul 29. Electronic tongue for pharmaceutical analytics: quantification of tastes and masking effects Legin A et al**

The organoleptic aspects of pharmaceutical formulations affect their acceptability to the patient and hence can have an important effect on concordance with treatment. Objective evaluation of these aspects, particularly the taste of the formulation and the drug substance it contains, is difficult. Whilst volunteer taste panels can be used to good effect their utility is limited, particularly during very early stage development when the toxicological profile of the active pharmaceutical ingredient (API) is yet to be established in detail. A potentiometric "electronic tongue" has

been applied to analyse a variety of 41 individual substances and mixtures of particular interest for pharmaceutical research and development. The electronic tongue (ET) was capable of discriminating between substances with different taste modalities and could also distinguish different substances eliciting the same basic taste; the ET is promising in terms of quantifying the content of each substance and has an ability to detect nuances of the basic taste (e.g. lingering or short-lived). After calibration the electronic tongue was successfully applied to predicting bitterness strength of binary mixtures with a sweetener in terms of "apparent" or "perceived" quinine content. In order to render a formulation palatable it is often necessary to mask the (usually bitter) taste of the API by the addition of masking agents such as sweeteners and flavours. The ET proved capable of distinguishing between formulations with different levels of sweetener and/or flavour in a manner that was consistent with their masking efficiency as perceived by a small human taste panel. A suitably calibrated ET could have the benefit of providing the pharmaceutical formulator with reliable data concerning the taste of the product quickly and with a reduced need to ask volunteers to taste active pharmaceutical samples. Early development activities could be facilitated when human tasting is usually not possible in the absence of the required toxicological data.

Thus if $m \geq n_0(\varepsilon) = max\{n_1, n_2\}$, we have

$$\sum_{l=1}^{n} \left| \frac{d}{dt} x_1(t + \tau_k) - \frac{d}{dt} x_1 ( (t + \tau_m) \right| < \varepsilon \qquad (2.21)$$

Which shows that

$$\lim_{k\to\infty} \frac{d}{dt} x_1(t + \tau_k) \quad exists\ for\ i = 1, 2, \dots, n$$

Uniformly on all compact subsets of $(-\infty, \infty)$. Thus we have

$$\lim_{k\to\infty} \lim_{k\to\infty} \frac{x_i(t + \tau_k + h) - x_i(t + \tau_k)}{h}$$

$$= \lim_{h\to 0} \lim_{k\to 0} \frac{x_l(t + \tau_k + h) - x_l(t + \tau_k)}{h} \qquad (2.22)$$

$$= \lim_{h\to 0} \frac{w_1(t+h) - w_1(t)}{h}, \quad i = 1, 2, \dots, n,$$

Showing that (since the left side of (2.22) exists) $w_i (i = 1, 2, \dots, n)$ is differentiable. Furthermore we have from (2.22) that

$$\frac{dw_l(t)}{dt} = \lim_{k\to\infty} [\, f_i \left(t + \tau_k, x_1(t + \tau_k), \dots, x_n\ (t + \tau_k)\right)$$

$$-f_l\big(t + \tau_k, w_1(t), \dots, w_n(t)\big)$$

$$+ - f_l\big(t + \tau_k, w_1(t), \dots, w_n(t)\big)]$$

$$= -f_l\big(t + \tau_k, w_1(t), \dots, w_n(t)\big); \quad i = 1, 2, \dots, n, \qquad (2.23)$$

By the uniform almost periodicity of $f_1, f_2, \dots, f_n$ on $R \times S$ and hence $w(t) = \big(w_1(t), \dots, w_n(t)\big)$is a solution of (1.1).

In order to complete the proof, we have to show that $w(t) = \big(w_1(t), \dots, w_n(t)\big)$ is an almost- periodic function whose module is contained in that of $\{f_1(t, x_1, \dots, x_n), f_2(t, x_1, \dots, x_n), \dots, f_n(t, x_1, \dots, x_n)\}$ for $(x_1, x_2 \dots, x_n) \in S$. For this purpose it is in fact sufficient (by the module containment theorem) to show that for any sequence of real numbers $\{\tau_k\}$ for which

$$\{f_1(t + \tau_k, x_1, \dots, x_n), f_2(t, + \tau_k, x_1, \dots, x_n), \dots, f_n(t, + \tau_k, x_1, \dots, x_n)\}$$

Converges uniformly on R for all $(x_1, \dots, x_n) \in S$, the sequence $\{w_1\ (t, + \tau_k)\ \dots, w_n(t + \tau_k)\}$ converges uniformly on R where $\tau_k$ tends to a finite number or $\pm\infty$ as $k \to \infty$.

In the case where $\tau_k \to \tau \neq \pm\infty$ as $k \to \infty$, we have $w_l\ (t, + \tau_k) \to w_i(t + \tau), i = 1, 2, \dots, n$ uniformly on R as $k \to \infty$. Let us suppose $\tau_k \to \infty$ as $k \to \infty$. For any $\varepsilon > 0$ there exists an integer $n_0(\varepsilon)$ such that if $m$ ☒ $n_0(\varepsilon)$ then

$$\sum_{l=1}^{n} |b_l\ (t + \tau_k) - b_l(t + \tau_m)| < (\varepsilon/4x^*)x_*, \qquad (2.24)$$

$$x^* \sum_{i=1}^{n}\sum_{j=1}^{n} |a_{ij}(t+\tau_k) - a_{ij}(t+\tau_m)| < \left(\frac{\varepsilon}{4x^*}\right)\mu x_* \qquad (2.25)$$

$$n(2x^*/x_*)\exp[\mu x_* \tau_k] < (\varepsilon/2x^*). \qquad (2.26)$$

Consider the function $v\big(w(s), w(s+\tau_m-\tau_k)\big)$ for $s \in [t, t+\tau_k]$ defined by

$$v\big(w(s), w(s+\tau_m-\tau_k)\big) = \sum_{l=1}^{n} |\log w_l(s) - \log w_l(s+\tau_m-\tau_k)|. \qquad (2.27)$$

Using the fact $w: \boldsymbol{R} \to S$ and $w$ is a solution of (1.1), calculate the right derivative $D^+v \ of \ v$;

$$D^+v\big(w(s), w(s+\tau_m-\tau_k)\big) = \sum_{i=1}^{n}\{\text{sign}[w_l(s) - w_i(s+\tau_m-\tau_k)]\}.$$

$$\left\{\frac{dw_i(s)}{ds}\cdot\frac{1}{w_l(s)} - \frac{dw_l(s+\tau_m-\tau_k)}{ds}\,\frac{1}{w_l(s+\tau_m-\tau_k)}\right\}$$

$$-\mu\sum_{l=1}^{n}|w_l(s) - w_l(s+\tau_m-\tau_k)| + \sum_{l=1}^{n}|b_i(s) - b_l(s+\tau_m-\tau_k)|$$

$$+x^*\sum_{l=1}^{n}\sum_{J=1}^{n}\left|a_{lJ}(s) - a_{lJ}(s+\tau_m-\tau_k)\right| \qquad (2.28)$$

$$-\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+\sum_{i=1}^{n}|b_l(s) - b_l(s+\tau_m-\tau_k)| \qquad (2.29)$$

$$+x^*\sum_{l=1}^{n}\sum_{j=1}^{n}\left|a_{lj}(s) - a_{ij}(s+\tau_m-\tau_k)\right|.$$

If we set $s-\tau_k = \sigma$ then we have from (2.29),

$$D^{+}v\big(w(s),w(s+\tau_m-\tau_k)\big)$$

$$\leq -\mu x_{*}v\big(w(s),w(s+\tau_m-\tau_k)\big)$$

$$+\sum_{l=1}^{n}|b_l(\sigma+\tau_k)-b_l(s+\ \tau_m)|$$

$$+x^{*}\sum_{l=1}^{n}\sum_{j=1}^{n}\left|a_{ij}\,(\sigma+\tau_k)-a_{ij}\,(\sigma+\tau_m)\right| \qquad (2.30)$$

$$\leq -\mu x_{*}v\big(w(s),w(s+\tau_m-\tau_k)\big)$$

$$+[(\varepsilon\mu x_{*})/(4x^{*})]\ +[(\varepsilon\mu x_{*})/(4x^{*})]$$

(2.31)

Integrating (2.31) over $(t,t+T_k)$

$$v(w(t+T_k)w(t+T_m))$$

Another MODEL where the model is directly applicable with the scheduled ones are the case of insurance and corporate finance. Insurance: Mathematics and Economics Volume 42, Issue 1, February 2008, Retrieval of Black–Scholes and generalized Erlang models by perturbed observations at a fixed time Dedicated to Professor André Dubey on the occasion of his 60th birthday Daniel Neuenschwander

S-stable laws on the real line (more generally on Hilbert spaces), associated with some non-linear transformations (so-called “shrinking operations”), were introduced in [Jurek, Z.J., 1977. Limit distributions for truncated random variables. In: Proc. 2nd Vilnius Conference on Probability and Statistics, June 28–July 3, 1977. In: Abstracts of Communications, vol. 3, pp. 95–96; Jurek, Z.J., 1979. Properties of s-stable distribution functions. Bull. Acad. Polon. Sci. Sér. Math. XXVII (1), 135–141; Jurek, Z.J., 1981. Limit distributions for sums of shrunken random variables. Dissertationes Math. vol. CLXXXV]. In [Jurek, Z.J., Neuenschwander, D., 1999. S-stable laws in insurance and finance and generalization to nilpotent Lie groups. J. Theoret. Probab. 12 (4), 1089–1107], the authors interpreted s-stable motions on the real line as limits of total amount of claims processes (up to a deterministic

premium) of a portfolio of excess-of-loss reinsurance contracts and showed that they led to Erlang's model or to Brownian motion. In [Neuenschwander, D., 2000b. On option pricing in models driven by iterated integrals of Brownian motion. In: Mitt. SAV 2000, Heft 1, pp. 35–39], we considered stochastic integrals whose integrand and integrator are both independent Brownian motions, thus modelling a stochastic volatility; as a result we got an analogue of the Black–Scholes formula in this model, confirming a result of Hull and White [Hull, J., White, A., 1987. The pricing of options on assets with stochastic volatility. J. Finance XLII (2), 281–300]. In the present paper, we will look at a common generalization of these processes, namely s-stable motions on the real line perturbed by a stochastic integral whose integrand and integrator are both (not necessarily independent) s-stable motions. The main result will be that if we can observe the distribution of such so-perturbed s-stable motions (together with the values of the perturbing processes) at time t=1, then we can identify the whole model (including the perturbation) among all models with Lévy processes perturbed by an iterated stochastic integral of two Lévy processes (in the Gaussian case) resp. among all models with a compound Poisson process with drift perturbed by an iterated stochastic integral of two compound Poisson processes (in the completely non-Gaussian case if the perturbing processes have no drift) without knowing anything about the history or about its distribution during 0≤t<1. This applies, e.g., to a situation where several assets obey the same model and one can estimate the distribution at time one by looking at the values of all these assets at time t=1. Interestingly enough, it will be convenient to treat the whole matter in the algebraic framework of the so-called Heisenberg group. This is a concept coming in fact from quantum mechanics and is in a certain sense the simplest non-commutative Lie group.

$ e^{-\mu x * T_k}\, v\big(w(t), w(t+T_m-T_k)\big) + (\varepsilon/2x^*)$

$ e^{-\mu x * T_k} \sum_{i=1}^{n} |\log\ w_i(t) - \log\ w_i(t+T_m-T_k)| + (1/\mathrm{x}_*) + (\varepsilon/2x^*)$

n(2x^* / x_*)exp[−μx_* T_k] + (ε/2x^*)

<(ε/x^*).

Since

$$v\big(w(t+T_k)w(t+T_m)\big) \sum_{i=1}^{n} |\ w_1(t+T_k) - w_1(t+T_m))/\ x^*)$$

We have from(2.33)that

$$\sum_{i=1}^{n} | w_1(t+T_k) - w_1(t+T_m)| < \varepsilon \; for \; m \; k \; h_0(\varepsilon)$$

Which show that the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\}$converges uniformly on **R**. By a similar argument the above conclusion is also valid when $T_k \to -\infty \; as \; k \to \infty$. Thus the uniform convergence of the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\} \; as \; k \to \infty \; on$ **R** follows and this completes the proof of the theorem.

**COROLLARY.** Suppose $b_i a_{ij}$ (i, j = 1.2, … , n)in (1.1)are periodic with a common period say $T >$ $0. Then \; if$ (1.2), (1.3)$and$ (2.3)$hold \; the \; system$ (1.1)$will \; have \; a \; periodic \; solution \; w(t) =$ $\{w, (t), \ldots\ldots, w_n(t)\} of \; period \; T \; such \; that$

$$w_1(t+T) = w_1(t), \qquad t \in \boldsymbol{R}: i = 1,2, \ldots, n.$$

**PROOF.** Proof follows from that of Theorem 2.1 if we Choose $T_k = KT(K = 0,1,2, \ldots) and \; note \; that$

$$x_1(t+KT) - w_1(t) \to 0 \; as \; k \to \infty; i = 1,2, \ldots, n; t \in \boldsymbol{R},$$

$$x_i(t+(K+1)t) - w_1(t) \to 0 \; as \; \boldsymbol{K} \to \infty; i = 1,2, \ldots, n; t \in \boldsymbol{R},$$

$$x_i(t+T+KT) - w_1(t+T) \to 0 \; as \; k \to \infty; i = 1,2, \ldots, n; t \in \boldsymbol{R},$$

Implying that

$$w_1(t+T) = w_i(t); \quad i = 1,2, \ldots, n; t \in \boldsymbol{R},$$

## 3. Global asymptotic stability

By definition we say that an almost-periodic solution say $w(t) = \{w_1(t), \ldots, w_n(t)\} of (1.1) is \; globally \; asymptotically$ stable (or attractive) if and only if every other solution x(t)={ $x_1(t), \ldots, x_n(t)$}with $x_1(o) > 0, i = 1.2, \ldots, n$ is defined for all t 0 and satisfies

$$\lim_{t\to\infty} |w_1(t) - x_1(t)| = 0; \quad i = 1,2, \ldots, n.$$

A consequence of such a global asymptotic stability is that there cannot be another strictly positive almost-periodic solution of (1.1)

**THEOREM 3.1** Suppose the almost- Periodic coefficients $b_1$, $a_{ij}$ (i, j = 1,2, … , n)

Satisfy (1.2)and (1.3)and the following

$$\inf_{t\in R} a_{ij}(t) > \left\{ \sum_{\substack{J=1 \\ J\neq 1}}^{n} \sup a_{ij}(t) \right\} + \mu; \qquad i = 1,2, \dots, n$$

For some positive number $\mu$ *Then* (1.1)has an almost-periodic solution (with strictly positive components) which is globally asymptotically stable.

We will omit the proof since the proof is identical with that of Theorem 3.1 of [3]

Under condition (3.2) the almost –periodic solution $w(t) = \{w_1(t), \dots . w_n(t)\}$ of (1.1)enjoys a somewhat strong from of stability even if the coefficients of (1.1)are subjected to a class of perturbations We note (see Yoshizawa [7]page17)that if a function b:R→ $R$ is almost periodic then H(b),the "hull of b" is defined as follows:

H(b)= {g|g:**R**→ $\boldsymbol{R}$ *and* $g(t)\lim_{k\to\infty} b(t + T_k)$ *for* some sequence

$T_k$ of real numbers the limit being uniform in $t \in R$.

By Theorem 2.2 of Yoshizawa [7, page 10], members of the hull of almost periodic functions are themselves almost periodic.

A system of the from

$$\frac{dy_1(t)}{dt} = y_1(t)\left\{B_i(t) - \sum_{J-1}^{n} A_{ij}(t) y_1(t)\right\},$$

*i=1,2,…,n; t>*$t_0 \in R$,

Where

$$B_i \in H(b_i); (A_{ij}), \qquad \text{i, j} = 1,2, \dots, \text{n},$$

is known as a system from the hull of (1.1) As a consequence of the definition of the hull in (3.3)and (3.5)we shall have

$$\inf_{t\in R} B_i(t) = \text{b}_i^{\,l}; \quad \sup_{t\in R} B_i(t) = \text{b}_i^{\,l};$$

$$\inf_{t\in R} A_{ij}(t) = \text{b}_i^{\,l}; \quad \sup_{t\in R} A_{ij}(t) = \text{a}_{iJ}^{\,u};$$

and hence the coefficients of(3.4)- (3.5)satisfy the same conditions as (1.2),(1.3)and (3.2)It will then follow from theorems(2.1)and (3.1)that solutions of (3.4)-(3.5)with $y_1(t_0) > 0, i = 1,2, \dots, n$ *will have* the following property:

$$y_1(t) > 0 \ \ for\ t \geq t_0; i = 1,2, \dots, n,$$

$$\sum_{i=1}^{n} y_1(t) remains\ bounded\ for\ t \geq t_0$$

**DEFINITION** A Solution $x(t) = \{x_1(t), \dots, x_n(t)\}, t \in t_0 \in \boldsymbol{R}$,of (1.1)is said to be stable under perturbations from the hull of (1.1)if the following holds :let $y(t) = \{(y_1(t), \dots, y_n(t)\}$ be any solution of (3.4)-(3.5)such that $y_1(t_0) > 0$, i=1,2,…,n,and for some positive constants $\beta_1, \beta_2 \beta_1, \leq y_1(t) \leq \beta_2, t \geq t_0, \quad i = 1,2, \dots, n$ Then given any $\varepsilon > 0$,there exists a constant $\delta > 0$ such that

$$\sum_{j=1}^{n} |x_1(t_0) - y_1(t_0)| < \delta$$

$$\sum_{i=1}^{n} \sup_{t\in R} |b_1(1) - B_1(t)| < \delta \qquad = \sum_{i=1}^{n} |x_1(t) - y_i(t)| < \varepsilon$$

$$\sum_{i=1}^{n}\sum_{j=1}^{n} \left| \sup_{t\in R} |a_{ij}(t) - A_{ij}(t)| < \delta \right.$$

For $t \geq t_0$.

We can now formulate the following result which shows that all solution of (1.1)are stable under perturbations from the hull of (1.1)

**THEOREM 3.2**. Assume that the conditions of Theorem (3.1) hold. Then every strictly positive solution of (1.1) ( including its unique almost-periodic solution) is stable under perturbations from the hull of (1.1).

PROOF. Let $x(t) = \{x_1(t), \dots, x_n(t)\}$ and $y(t) = \{y_1(t), \dots, y_n(t)\}$ be any two strictly positive solutions of (1.1) and (3.4) – (3.5) respectively for $t \geq t_0 \in R$. Let $\beta_1, \beta_2$ be constants such that

$$\begin{matrix} \beta_1 \leq x_i(t) \leq \beta_2, \\ \beta_1 \leq y_i(t) \leq \beta_2, \end{matrix} \quad i = 1,2,\dots,n; t \geq t_0 \qquad (3.7)$$

Consider now a function $V(t) = V(x(t)y(t))$ defined by

$$V(x(t)y(t)) = \sum_{i=1}^{n} \left| \log x_1(t) - \log, y_i(t) \right|, \; t \geq t_0 \in R. \qquad (3.8)$$

It is easy to see that

$$\frac{1}{\beta_2} \sum_{i=1}^{n} \left| x_i(t) - y_i(t) \right| \leq V(x(t)y(t))$$

$$\leq \frac{1}{\beta_1} \sum_{i=1}^{n} \left| x_i(t) - y_i(t) \right|, \; t \geq t_0 \in R. \qquad (3.9)$$

Calculating the right derivative $D^+V$ of $V$ and using (1.1) and (3.4) we derive (after simplification) that

$$D^+V\left(x(t),y(t)\right) \le -\mu\sum_{i=1}^{n}\left|x_i(t)-y_i(t)\right| + \sum_{i=1}^{n}\left|b_i(t)-B_i(t)\right|$$

$$+\beta_2\sum_{j=1}^{n}\sum_{i=1}^{n}\left|a_{ij}(t)-A_{ij}(t)\right|; \quad t \ge t_0,$$

$$\le -\mu\,\beta_1 V\left(x(t),y(t)\right) + \sum_{i=1}^{n}\left|b_i(t)-B_i(t)\right|$$

$$+\beta_2\sum_{i=1}^{n}\sum_{j=1}^{n}\left|a_{ij}(t)-A_{ij}(t)\right|; \quad t \ge t_0, \tag{3.10}$$

An integration of (3.9) over $[t_0, t]$ with an application of (3.8) leads to

$$\frac{1}{\beta_2}\sum_{i=1}^{n}\left|x_i(t)-y_i(t)\right| \le V(x(t),y(t)) \le V(x(t_0),y(t_0))$$

$$+\left[\sum_{i=1}^{n}\sup_{t\in R}\left|b_i(t)-B_i(t)\right| + \beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}\left|a_{ij}(t)-A_{ij}(t)\right|\right]\frac{1}{\mu\,\beta_1},$$

$$t \ge t_0,$$

And hence

$$\sum_{i=1}^{n}\left|x_1(t)-y_1(t)\right| \le \left(\frac{\beta_2}{\beta_1}\right)\sum_{i=1}^{n}\left|x_1(t_0)-y_1(t_0)\right|$$

$$+\frac{\beta_2}{\mu\beta_2}\left[\sum_{i=1}^{n}\sup_{t\in R}\left|b_i(t)-B_i(t)\right|\right.$$

The construction reason is because is criticized in a number of usage guides: The reason for the long delays was because the costs greatly exceeded the original estimates. One objection to this construction is based on its redundancy: the word because (literally, by cause) contains within it the meaning of reason; thus saying the reason is because is like saying "The cause is by cause," which would never be said. A second objection is based on the claim that because can introduce only adverbial clauses and that reason is requires completion by a noun clause. Critics would substitute that for because in the offending construction: The reason for the long delays in completing the project was that the costs.

$$+\beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}\left|\ a_{ij}(t)-A_{ij}(t)\right|\Big]. \qquad (3.11)$$

Now for any $\varepsilon>0$ if one chooses a $\delta>0$ such that

$$\delta<\min\left\{\frac{\varepsilon}{3}\ \frac{\beta_1}{\beta_2},\frac{\varepsilon}{3}\ \frac{\beta_1}{\beta^2}\mu\right\} \qquad (3.12)$$

Then (3..6) will be satisfied and this completes the proof

## . Jia Li's Presciential Dynamics And Contential Staticity Of Accentuation Attrition Models Applied To Cosmological And Astrophysical And Quantum Mechanical Engineering Problems

**Abstract**: **Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics ET AL ). We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept**

**formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems. Predator prey model literally form bastion, pillar, post, stylobate, sentinel towards the end of dismasting many of the asymptotic and solutional behaviour in finance, economic, and as we shall see in cosmology, quantum computation and other fields.**

## MODEL FORMULATION

Let $p_1(t, a_1)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_1}{dt} = -b_1 Q_1 + g_1(g_1)Q_1$$

Where $b$ is the natural death rate in the absence of credit types and $g(P)$ is the conversion rate function of debits

into credit creations which is assumed to be differentiable and satisfies

$$g_1(0) = 0, \quad g_1(P_1) > 0, \ P_1 \ 1, \qquad \lim_{P_1 \to \infty} g_1(P_1) = (K_1)_1 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_1 + (\rho_a)_1 + D_1(a_1, P_1, Q_1)\rho_1 + 0, \qquad 0 < a_1 \ \infty,$$

$$\rho_1(t,0) = \int_0^\infty n_1(a_1, P_1, Q_1)\rho_1(t_1, a_1)da_1$$

Where $D_1$ is the mortality function and $n_1$ is the fertility function.

Put $_1$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_1}{dt}\right]_1 = x_1\left[\frac{dQ_2}{dt}\right]_1 - x_2\left[\frac{dQ_1}{dt}\right]_1$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_1 - x_4\left[\frac{dQ_2}{dt}\right]_1$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_1 - x_6\left[\frac{dQ_3}{dt}\right]_1$$

Where $x_1, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_1$ and $g_1(P_1)Q_1$

Now when you add the term $(b_1Q_1)$ with $g_1(P_1)(Q_1)$ we are infect see is that $(b_1Q_1)$ is credited incrementally part of $g_1(P_1)(Q_1)$ or instant nearly, when subtracted from $\frac{dQ_1}{dt}$ , it means that $g_1(P_1)(Q_1)$ is slowly eaten by $\frac{dQ_1}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_1 = \mu_1(a_1) + (f_1)_1(P_1)_1 + (f_2)_1(Q)_1 \quad (2)$$

$$n_1 = \beta_1(a_1)h_1(Q_1), \quad (3)$$

Where $\mu_1$the age- dependent debit rate whit is $\mu_1(a_1)$ 0 $\forall a > 0$; $(f_1)_1$ is the density- dependent death rate with $f_1(0) = 0$, $(f_1^1)_1(P_1)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_1)_1(P_1) = (k_1)_1$ ; $f_2$ is the functional response of creditors

to prey with $[f_2]_1(0) = 0, (f_2^1)_1(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_1\ (Q_1) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^1(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_1 + (\rho_a)_1 = -\mu_1(a_1) + (f_1)_1 P_1) + (f_2)_1(Q_1))(\rho_1)_1 \quad (4)$$

$$\rho_1(t_1, 0) = h_1(Q_1)\int_0^\infty \beta_1(a_1)\rho_1(t, a_1)da, \quad (5)$$

$$\frac{dQ_1}{dt} = -b_1 Q_1 + g_1(P_1)Q_1 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_1$ and $Q_1$ of (4-6) satisfy

$$0 \ \rho_1(t, a_1) \ \rho_1^*(t, a_1),$$

$$0 \ Q_1(t) \ Q_1^*(t),$$

Where $\rho_1^*(t, a_1)$ and $Q_1^*(t)$, are solutions of

$$(\rho_t^*)_1 + (\rho_a^*)_1 = -(\mu_1(a_1) + (f_1)_1(P_1^*) + (f_2)_1 Q_1^*))\ \rho_1^*, \quad (7)$$

$$\rho_1^*(t, 0) = \int_0^\infty \beta_1(a_1)\ \rho_1^*(t, a_1)da_1 \quad (8)$$

$$\frac{dQ_1^*}{dt} = -b_1\ Q_1^* + g_1(P_1^*)Q_1^*, \quad (9)$$

With $P_1^*(t) = \int_0^\infty \rho_1^*\ (t, a_1)da_1$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_1^*(t)}{dt} = P_1^*(t)\big(r_1 - (f_1)(P_1^*) - (f_2)_1(Q^*)\big), \quad (10)$$

$$\frac{dQ_1^*(t)}{dt} = Q_1^*(t)(-b_1 + g_1(P_1^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_1(a)e^{-\int_0^a \mu_1(\tau)dt} e^{-ra}\ da_1 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_1^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(t)dt}\ da_1 > 1, \quad (H1)$$

Which guarantees $r_1 > 0$.

*Remark.* If $\beta_1(a) = [\beta_0]_1 a_1 e^{-\alpha_1 a_1}, \mu_1(a_1) \equiv \mu_1$ then the solution of (12) is

$$r_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_1]_1(P_1) = 0, [f_2](Q_1) = \lambda_1 Q_1\ h_1(Q_1) = 1, \qquad and\ \ g_1(P_1) = c_1 P_1$$

$$[\rho_1]_1 + [\rho_a]_1 = -(\mu_1 + \lambda_1 Q_1)\rho_1$$

$$\rho_1(t,0) = \int_0^{\infty} [\beta_0]_1 a_1 e^{-\alpha_1 a_1} \rho_1(t,a_1)\, da_1$$

$$\frac{dQ_1}{dt} = -b_1 Q_1 + c_1 P_1 Q_1$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_1}{dt} = \omega_1 P_1 - \lambda_1 P_1 Q_1$$

$$\frac{dQ_1}{dt} = b_1 Q_1 - c_1 P_1 Q_1$$

Where $\omega_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$ THE EQUILIBRIUM $(\rho_0)_1(a_1), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_1(a_1), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_1 + (\rho_a)_1 = -\big(\mu_1(a_1) + (f_1)_1(P_1)\big)\rho,$$

$$\rho_1(t,0) = \int_0^{\infty} \beta_1(a_1)\rho_1(t,a_1)da_1 \qquad (13)$$

The solution of (13) has the form of

$$\rho_1(t,a_1) = Ce_1^{r(t-a_1)} - \int_0^{a} (\tau)d\tau - \int_0^{1} (f_1)_1\big(P_1(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_1(t) = C_1 e^{r_1 t - \int_0^t (f_1)_1(P_1(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_1^{-1}(r)$ if $k_1 > r$.

This is equivalent to

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{(k_1)_1 a_1} da_1 < 1.$$

In fact, if we define

$$R_1(s_1) = \int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau)d\tau} e^{-s_1 a_1}\, da_1$$

$R(s)$ is a decreasing function. $R_1(r_1) = 1$ And $R_1(k_1)_1 < 1$ imply $r_1 < (k_1)_1$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_1(a_1), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{k_1 a_1} da_1 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_1 \equiv \int_0^\infty (\rho_0)_{11}(a_1)da_1$ *satisfies*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{-(f_1)_1(P_0)a_1} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_1(a_1), 0)$

Let $u_1(t, a_1) = \rho(t, a_1) - \rho_0(a_1)$ Then

$$N_1(t) \equiv \int_0^\infty u_1\,(t, a_1)da_1 = \int_0^\infty \rho_1(t, a_1)da_1 - \int_0^\infty (\rho_0)_1(a_1)da = P_1(t) - P_0)_1$$

And

$$(u_1)_1 + (u_a)_1 = (\rho_t)_1 + (\rho_a)_1 - (\rho_0^1)_1(a_1)$$

$$= -\mu_1(a_1) + (f_1)_1(P_1) + (f_2)_1(Q_1))\,(u_1 + (\rho_0)_1) + \mu_1(a_1)(\rho_0)_1 + (f_1)_1(P_0)_1(\rho_0)_1$$

$$= -\mu_1(a_1)u_1 - (f_1)_1(P_1) + (f_2)_1(Q_1))\,u_1 - (\rho_0)_1\,(f_1)_1(P_1) - (f_1)_1(P_0)_1$$

$$+ (f_2)_1(Q_1)), \qquad (15)$$

Where $(\rho_0^1)_1(a_1)$ denotes $d(\rho_0)_1(a_1)/da_1$

Expansion $(f_1)_1(P_1)$ and $(f_2)_1(Q_1)$ at $(P_0)_1$ and 0 yield respectively

$$(f_1)_1(P_1) = (f_1)_1((P_0)_1 + N_1)$$

$$= (f_1)_1(P_0)_1 + (f_1^1)_1(P_0)_1N_1 + o_1(N_1), \qquad (16)$$

$$(f_2)_1(Q_1)$$

$$= (f_2^1)_1(0)Q + o_1(Q_1) \qquad (17)$$

Write $f_1(P_0) = f_{10}, f_1^{'}(P_0) = f_{10}^{'}$, and $f_2^{'}(0) = f_{20}^{'}$ and substitute (16) and (17) into (15) to obtain

$$(u_1)_1 + (u_a)_1 \approx -(\mu_1(a_1) + (f_{10})_1)u - (f_{10}^1)_1(\rho_0)_1(N_1) - (f_{20}^1)_1(\rho_0)_1(Q_1)$$

$$u_1(t, 0) = \rho_1(t, 0)(\rho_0)_1(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t, a) + (\rho_0)_1(a_1)da_1 - (\rho_0)_1\,(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t, a_1)da_1 + (h_1(Q_1) - 1)(\rho_0)_1\,(0) \qquad (18)$$

Expanding

$$h_1(Q_1) = 1 + (h^1)_1(0)Q_1 + 0(Q_1) \equiv 1 + h_0' Q_1 + 0Q_1$$

And substituting it into (18), we arrive at

$$u_1(t,0) \approx \int_0^\infty \beta_1(a_1)u_1(t,a_1)da_1 + (h_{a_1}')(\rho_0)_1 Q_1$$

The linearization is the following:

$$[u_1] + [u_1] = -(\mu_1(a_1) + [t_{10}]_1 u_1 - [\alpha_1]_1(a_1)N_1 - [\alpha_2]_1(a_1)Q_1$$

$$u_1(t,0) = \int_0^\infty \beta_1(a_1)u_1(t,a_1)da_1 + [\xi_1]Q_1 \qquad (19)$$

$$\frac{dQ_1}{dt} = -b_1 Q_1$$

where

$$[\alpha_1]_1(a_1) = [t_1^1 a]_1[p,o]_1(a_1)d2(a_1) = [f_{20}^1]_1(\rho\, o\,)_1(a_1)\xi = [\mathrm{h\ o}^1]_1[\rho\mathrm{o}]_1(\mathrm{o}) \text{ and } \mathrm{b}_1 = \mathrm{b}_1 - g_1(P_1)$$

If $g_1(P_0) > b$, clearly, $(\rho_0)_1(a_1), 0)$ is unstable and the predator population is persistent. Assume $g_1(P_0)_1 < b$.

Then, $Q_1(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_1 + u_a = -(\mu_1(a_1) + (f_{10})_1)u_1 - (\alpha_1)_1(a_1)N_1$$

$$= -(\mu_1(a_1) + f_{10})u_1 - \alpha_1(a_1)\int_0^\infty u_1(t,a_1)da_1$$

$$u_1(t,0) = \int_0^\infty \beta_1(a_1)u_1(t,a_1)da_1$$

Using the method developed by Cushing [2, 3], we suppose

$$u_1(t,a_1) = \tilde{u}_1(a_1)^{e^{\gamma(t-a_1)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_1}\, da_1$. Then,

A solution of the first differential equation for $\tilde{u}(a_1)$ gives,

$$\tilde{u}_1(a_1) = e^{-\int_0^a \mu_1(\tau_1)d\tau - (f_{10})_1 a}\left(\tilde{u}_1(0) - w\int_0^a e^{\int_0^\tau \mu_1(s_1)ds + (f_{10})}\tau_1\,(\alpha_1)_1(\tau_1)e^{\gamma_1\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_1(a_1)e^{-\gamma_1 a_1 - (f_{10})_1 a_1 - \int_0^a \mu_1(\tau_1)d\tau}\right)\tilde{u}_1(0)$$

$$+w\int_0^\infty \beta_1(a_1)e^{-\gamma_1 a_1 - (f_{10})_1 a_1 - \int_0^a \mu_1(\tau_1)d\tau}\int_0^a (\alpha_1)_1(\tau_1)e^{-\gamma_1\tau_1}e^{(f_{10})_1\tau_1 + \int_0^\tau \mu_1(s_1)ds}d\tau_1\, da_1$$

$$= 0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-ya - f_{10}a - \int_0^a \mu(\tau)d\tau}da\right)u(0)$$

$$+\left(-1 - \int_0^\infty e^{-ya - f_{10}a - \int_0^a \mu(\tau)d\tau}\int_0^a e^{f_{10}\tau + \int_0^\tau \mu(s)ds + y\tau}\alpha_1(\tau)d\tau\, da\right)w = 0. \quad (23)$$

Define $Y(\gamma,a) = e^{-ya - f_{10}a - \int_0^a \mu(\tau)d\tau}$. Then

$$\left(1 - \int_0^\infty \beta\,(a)Y(\gamma,a)da\right)u(0)$$

$+w\int_0^\infty \beta(a)Y(\gamma,a)\int_0^a \alpha_1(\tau)Y^1(\gamma,\tau)d\tau\, da = 0, (24)$

$$\left(\int_0^\infty Y(\gamma,a)da\right)u(0)$$

$$+\left(-1-\int_0^\infty Y(\gamma,a)\int_0^a \alpha_1(\tau)Y^{-1}(\gamma,\tau)d\tau\, da\right)w = 0.$$

Equation (20) has a solution of the form of $u(t,a) = u(a)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $u(0) \not\equiv 0$ and $\gamma, w$ in the complwx plane.

Let $I(y,a) = \int_0^a \rho_0(\tau)Y^{-1}(\gamma,\tau)d\tau$. We obtain the characteristic equation

$$f_{10}^l\left(\int_0^\infty Y(\gamma,a)da\right)\left(\int_0^\infty \beta(a)Y(\gamma,a)I(\gamma,a)da\right)$$

$$+\left(1-\int_0^\infty \beta(a)Y(\gamma,a)da\right)\left(1+f_{10}^l\int_0^\infty Y(\gamma,a)I(\gamma,a)da\right) = 0. \quad (25)$$

Since

$$\rho_0(a) = \rho_0(0)e^{-\int_0^a \mu(\tau)d\tau - f_{10}a}$$

$$I(\gamma,a) = \int_0^a \rho_0(\tau)e^{\gamma\tau + f_{10}\tau + \int_0^\tau \mu(s)ds}d\tau$$

$$= \frac{\rho_0(0)}{\gamma}(e^{\gamma a} - 1)$$

Let

$$E(a) = e^{-f_{10}a - \int_0^a \mu(\tau)d\tau}$$

Then $[Y_1]_1[\gamma_1]_1, a_1) = E_1(a_1)e^{-\gamma_1 a_1}$. by a simple calculation,

$$\int_0^\infty \beta_1(a_1)Y_1(\gamma_1, a_1)I(\gamma_1, a_1)da = \frac{[\rho_0]_1(0)}{\gamma}\left(1 - \int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}\ da\right);$$

$$\int_0^\infty Y_1(\gamma_1, a_1)I(\gamma_1, a_1)da = \frac{[\rho_0]_1(0)}{\gamma_1}\left(\int_0^\infty E_1(a_1)e^{-\gamma_1 a_1}\ da_1\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_1(0)}{\gamma_1}[f_{10}^1]\int_0^\infty E_1(a_1)e^{-\gamma_1 a_1}da_1\ (\int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1)$$

$$+\left(1 - \int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right).$$

$$\left(1 + \frac{(\rho_0)_1(0)}{\gamma_1}(f_{10}^{'})_1\left(\int_0^\infty E_1(\ a_1)da_1 - \int_0^\infty E_1(\ a_1)e^{-\gamma_1 a_1}da_1\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1 - \int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right)(1 + \frac{(f_{10}^{'})_1}{\gamma_1}(\rho_0)_1(0)\int_0^\infty E_1(\ a_1)da_1) = 0.$$

We need ũ(a) $\not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1 - \int_0^\infty \beta_1(a_1)Y_1(\gamma_1 a_1)da_1 = 1 - \int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1 \neq 0,$$

Since $\int_0^\infty \beta_1(a_1)Y_1(\gamma_1 a_1)\ \int_0^a \alpha_1\ (\tau_1)Y^{-1}(\gamma_1,\tau_1)d\tau_1 da_1 > 0.$ Therefore, the characteristic solution is

$$\gamma_1 = -(f_{10}^{'})_1(\rho_0)_1 \int_0^\infty E_1(a_1)da_1.$$

Since $f_1'(P_0)_1 > 0,$ we obtain the following result

THEOREM 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_1), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_1), 0)$ which is related to the extinction of debit summation as has been referred to in the foregoing and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(\ p^*(a), Q^*\ )$. It must satisfy the equations:

$$[p_a^*]_1 + (\mu_1(a_1) + [f_1]_1(\ P_1^*) + [f_2]_1(Q_1^*)\ )P_1^* = 0, \quad (26)$$

$$p_1^*(0) = h_1(Q_1^*)\int_0^\infty\ \beta_1(a_1)p_1^*(a_J)p_1^*(a_1)da_1 \quad (27)$$

$$g_1(P_1^*) - b_1 = 0, \quad (28)$$

Where $P_1^*$=$\int_0^\infty p_1^*(a_1)\, da_1$

Equation (28) yields $P^* = g^{-1}(b),$

From (26),

$$p_1^*(a) = p_1^*(0)e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}$$

Substituting into (27), we need to solve

$$\left(1 - h_1(Q_1^*)\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}da_1\right)p_1^*(0) = 0$$

For $Q_1^*$

Fix $p_1^*(0)$ and $P_1^*$ Let

$$H_1(z_1) = 1 - h_1(z_1)\int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1}da_{,1}$$

And assume that

$$[k_3]_1\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(g_1^{-1}(b_1))a_1 - [k_2]_1 a_1}\, da_1 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_1(z_1)] > 0.$$

Clearly,

$$\frac{dH_1}{dz_1} = h_1(z)([f_2']_1(z_1))\int_0^\infty a_1\beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - f_1(P_1^*)a_1 - f_2(z_1)a_1}da_{1}$$

$-h'(z_1)\int_0^\infty \beta_1(a_1)\, e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1}\quad da_1$

By the assumptions on $\left[f_2^1\right]_1 and\ [h^1]^1 dH_1/dz > 0$. Hence, if

$$H_1(0) = 1 - \int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1} da_1 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_1(0) < 0$ follows if $g_1^-(b_1) < [f_1^{-1}]_1(r_1)$, where $r_1$ is the solution of (12).

Indeed, since

$$R_1(s_1) = \int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 e^{-y}1^a} da_1$$

Is a decreasing function of $s_1$ when $P_1^* = [g^{-1}]_1(b_1) < [f_1^{-1}]_1(r_1)$, that is $[f_1]_1(P_1^*) < r$, $R\left(f_1(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_1(g^{-1}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1\left(g_1^{-1}(b)\right)a_1} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_1^*(a_1), Q_1^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_1^*(a_1), Q_1^*)$ is more complicated. We again consider

perturbations $m_1(t_1 a_1) = p_1(t_1, a_1) - p_1^*(a_1)$ and $W_1(t) = Q_1(t) - Q^*$.Then

$M_1(t) \equiv \int_0^\infty m_1(t, a_1) da_1 = P_1(t) - P_1^*,$

$$[m_1]_1 + [m_a]_1 = [p_1]_1 + [p_a]_1 - \left[p^*\right]_1(a_1) \approx -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1$$

$$-[f_{1*}^{1}]_1 p_1^*(a_1)M - [f_{2*}^{1}]_1 p_1^*(\mathrm{a})W_1,$$

$$m_1(t,0)p_1(t,0) - p_1^*(0) \approx h_* \int_0^\infty \beta_1(a_1)m_1(t,a_1)\,da + \frac{[h_*^1]_1}{[h_*]_1}\, p_1^*(0)W_1$$

$$\frac{dW_1}{dt} = \frac{dQ_1}{dt} \approx [g_*^1]_1 Q_1^* M_{,1}$$

Where $[f_{1*}]_1$ denoting $[f_1(P^*)]_1$, $[f_{1*}]_1$ denoting $[f_2(P^*)]_1$, $\left[f_{1*}^{'}\right]_1$ denoting $\left[f^{'}{}_1(P^*)\right]_1 \left[f_{2*}^{'}\right]_1$ denoting

$[f_2]_1\,(Q_1^*,\quad h_{1}^{*}$ denoting $h_*^{'}$ denoting $h_1^{'}\,(Q_1^*)$, and $\left[g_*^{'}\right]_1$

Denoting $\left[g^{'}\right]_1(P_1^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_1]_1 + [m_a]_1 = -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1 - \left[f^{'}{}_{1*}\right]_1 [p_1^*](a_1)M_1 - \left[f_{2*}^{'}\right]_1 p_1^*(a_1)W_1$$

$$m_1(t,0) = [h_*]_1 \int_0^\infty \beta_1(a_1)m_1(t,a_1)da_1 + \frac{[h_*^1]_1}{[h^*]_1}[p^*]_1(0)\,W_{,1} \qquad (30)$$

$$\frac{dW_1}{dt} = [g_*^1]_1 [Q_1^*]M_1$$

Integrating $[m_1]_1(t,a_1)$ with respect to a from 0 to ∞, we have

$$\frac{dM_1(t)}{dt} = \int_0^\infty [mt]_1(t,a_1)da_1$$

$$= -\int_0^\infty [m_a]_1\,(t,a_1)da_1 - \int_0^\infty \mu(a_1)m(t,a_1)da_1$$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f'_{1*}]_1 \int_0^\infty p_1^*(a_1)da_1\right) M_1 - [f_{2*}^1]_1 \int_0^\infty p_1^*(a_1)da_1 W_1$$

$= m_1(t,0) - \int_0^\infty \mu_1(a_1)\, m_1(t,a_1)da_1$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f_{1*}^{i}]_1 P_1^*\right) M_1 - [f_{2*}^{i}]_1 P_1^* W_{\cdot 1}$$

From (30), we arrive at the equations

$$\frac{dW_1(t)}{dt} = \int_0^\infty ([h_*]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1$$

$$-\left([f_{1*}] + [f_{2*}]_1 + [f_{1*}']_1 P_1^*\right) M_1 + \left(\frac{[h_*']_1}{[h_*]^1}\, P_1^*(0) - f_{2*}' P_1^*\right) W_1,$$

(31)

$$\frac{dW_1(t)}{dt}[g'_*]_1 Q_1^*\, M_{\cdot 1}$$

By the mean value theorem,

$$\int_0^\infty ([h_{"*}]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1 = ([h_*]_1\beta_1(\theta_1) - \mu_1(\theta_1))M_{,1}$$

Where $0< \theta_1 = \theta_1(t) < \infty$.

And if $a_1^2 [a_1^2]_1 < 4 [a_2]_1 < 4 a_2 e$,

Re $\lambda_1[\lambda_1]_1 = \mathrm{Re}\lambda_2[\lambda_2]_1 = -\frac{1}{2}a_1[a2\ ]_1 < 0.$

Let $\Delta$=max$_{2<u<\infty}\left|h_*[h_*]_1\beta_1(a_1) - \mu_1(a_1)\right|$. Then$|n_1(t)|$ ⊿∀t ⓪. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium($\rho^*(a)$, $\boldsymbol{Q}^*$) is asymptotically stable if either $a_1^2 < 4\,a_2 e$ and $[a_1] > 2\Delta$,or $[a_1^2]$ ④$[a_2]$e and $(a_1 - \sqrt{(a_1^2)_1 - 4(a_2)_1 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+1}\right)$ and $\beta_1(a_1) = ae^{-ma}$ Then $\Delta$=$[\mu_0] + 1$ Simple calculation leads to

$$\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(\tau_1) + d\tau_1 da} = \frac{1}{(m+\mu_0+1)^2} + \frac{2}{(m+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_1 + [\mu_0]$ ⓪.5. Assume$[f_1](x_1)$= $K_1x_1$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_1(x_1)$= $[0x_1]/(1+x_1)$; the functional responses $[[f_2](x) = [0]\,x_1/(1+x_1)$and $h_1(x_1)$=$(1+\omega_1x_1)/(1+x_1)$(See Holling [7] and Rashevsky [13])

Equation (14),now has the form of

$$\frac{m_1 + [\mu_0] + [\mathrm{K}][P_0])}{[m_1]_1 + (\mu_0)_1 + [\mathrm{K}][P_0])_1^3\ [\mathrm{K}]_1[P_0]_1^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1+P_0) < b_1$ i e $[P_0] < b/(\delta_1 - b_1)$then the equilibrium$(P_0)$ $[P_0]_1(a), 0)$ is stable This implies that the

cedit confederation goes to extinction but that the cedit confederation persists $b_1$

$[f_{1*}] = K_1$ P $^*$ $[f_{2*}] = \frac{\sigma Q_1^*}{1+Q_1^*}$ $[h_*] = \frac{1+\omega_1 Q_1^*}{1+Q_1^*}$

$[f_{1*}] = K_1$ $\left[f_{2*}^{1}\right] = \frac{\sigma_1}{(1+Q_1^*)^2}$ $[g_*^1] = \frac{\sigma_1}{(1+P_1^*)^2}$ $\left[h_*^1\right] = \frac{\omega_1 - 1}{(1+Q_1^*)^2}$

## AN ANALOGY BETWEEN PREDATOR PREY AND ACCENTUATION ATTRION MODELS

**Best example how the model could be applied is the Schrödinger's equation itself wherein RHS gobbles up LHS, and the terms on RHS do the same with each other. It is like adding milk to water. Whether it happens instantaneously or time lag is there is a different question altogether dealt at the opportune points**

**Best example that can be given is the formulation of the quantum wave equation albeit in a nascent form for the recently found gravitational waves. This has been studied in length in various papers of gravitational waves, wherein the attribution of how the model is used is described. Note the calculation holds for only one module. This is in consideration to the fact that the relation is beteen the both debit systems and credit systems which as we have enucleated in the foregoing bears ample testimony for the application of the model**

**Gravitational quantum wave equation:**

**Equation Time-dependent equation (Wikipedia)**

**(1) In quantum mechanics, the Schrödinger equation is (=) a partial differential equation that describes how the quantum state of a quantum system changes with (e&eb) time. It was formulated in late 1925, and published in 1926, by the Austrian physicist Erwin Schrödinger.[1]**

**(2) In classical mechanics Newton's second law, (F = ma), is used to make (eb) a mathematical prediction as to what path a given system will take following a set of known initial conditions.**

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| **(3) In quantum mechanics, the analogue of Newton's law is (=) Schrödinger's equation for (e) a quantum system (usually atoms, molecules, and subatomic particles whether free, bound, or localised).** |
| **(4) It is not a simple algebraic equation, but in general a linear partial differential equation, describing (eb) the time-evolution of the system's wave function (also called a "state function").[2]:1–2** |
| **(5) The concept of a wavefunction is (=) a fundamental postulate of quantum mechanics.** |
| **(6) Although Schrödinger's equation is often presented as a separate postulate, some authors [3]: Chapter 3 show that some properties resulting from Schrödinger's equation may be deduced just from (e) symmetry principles alone, for example the commutation relations.** |
| **(7) Generally, "derivations" of the Schrödinger equation demonstrate (eb) its mathematical plausibility for describing (eb) wave–particle duality, but to date there are no universally accepted derivations of (e) Schrödinger's equation from (e) appropriate axioms.** |
| **(8) In the Copenhagen interpretation of quantum mechanics, the wave function is (=) the most complete description that can be given of (e) a physical system.** |
| **(9) Solutions to Schrödinger's equation describe (eb) not only molecular, atomic, and subatomic systems, but also macroscopic systems, possibly even the whole universe.[4]:292ff** |
| **(10) The Schrödinger equation is consistent with (e&eb) classical mechanics but not with (e&eb) special relativity.** |
| **(11) Making quantum mechanics consistent with (e&eb) special relativity is one of the great achievements of (e) Quantum Field Theory and requires (e) the introduction of the creation and annihilation operators so that (e) the number of particles can no longer be considered constant.** |
| **(12) The Schrödinger equation is not the only way to make predictions in (eb) quantum mechanics—other formulations can be used, such as Werner Heisenberg's matrix mechanics, and Richard Feynman's path integral formulation** |
| **(13) The form of the Schrödinger equation depends on (=) the physical situation (see below for special cases).** |
| **(14) The most general form is the time-dependent Schrödinger equation, which gives a description of a system evolving with time: [5]:143** |

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| **(15) A wave function that satisfies the non-relativistic Schrödinger equation with V = 0. In other words, this corresponds to (e &eb) a particle traveling freely through (e&eb) empty space. Time-dependent** |
| **Schrödinger equation (general)** |
| **(16) i ℏ ∂ ∂ t Ψ ( r , t ) = H ^ Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t) where i is the imaginary unit, ħ is the Planck constant divided by 2π, the symbol ∂/∂t indicates a partial derivative with respect to time t, Ψ (the Greek letter psi) is the wave function of (e) the quantum system, and Ĥ is the Hamiltonian operator (which characterizes the total energy of any given wave function and takes different forms depending on the situation).** |
| **(17) Each of these three rows is a wave function which satisfies (eb) the time-dependent Schrödinger equation for a harmonic oscillator. Left: The real part (blue) and imaginary part (red) of the wave function. Right:** |
| **(18) The probability distribution of finding the particle with (e&eb) this wave function at a given position.** |
| **(19) The top two rows are examples of stationary states, which correspond to (e&eb) standing waves.** |
| **(20) The bottom row is an example of a state which is (=) not a stationary state.** |
| **(21) The right column illustrates (eb) why stationary states are (=) called "stationary".** |
| **(22) The most famous example is the non-relativistic Schrödinger equation for (e) a single particle moving in an electric field (but not a magnetic field; see the Pauli equation): [6]** |
| **Time-dependent Schrödinger equation** |
| **(Single non-relativistic particle)** |
| **(23) i ℏ ∂ ∂ t Ψ ( r , t ) = [ − ℏ 2 2 μ ∇ 2 + V ( r , t ) ] Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)** |
| **(24) Here μ is (=) the particle's "reduced mass", V is its potential energy, ∇2 is the Laplacian (a differential operator), and Ψ is (=) the wave function (more precisely, in this context, it is called the (=) "position-space wave function").** |

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| **(25) In plain language, it means "total energy equals (=) kinetic energy plus potential energy", but the terms take unfamiliar forms for (e) reasons explained below.** |
| **(26) Given the particular differential operators involved, this is (=) a linear partial differential equation.** |
| **(27) It is also a diffusion equation, but unlike the heat equation, this one is (=) also a wave equation given (e&eb) the imaginary unit present in the transient term.** |
| **(28) The term "Schrödinger equation" can refer to (e) both the general equation (first box above), or the specific nonrelativistic version (second box above and variations thereof).** |
| **(29) The general equation is indeed quite general, used throughout quantum mechanics, for everything from the Dirac equation to quantum field theory, by plugging in various complicated expressions for the Hamiltonian. The specific nonrelativistic version is a simplified approximation to reality, which is quite accurate in many situations, but very inaccurate in others (see relativistic quantum mechanics and relativistic quantum field theory)** |
| **(30) To apply the Schrödinger equation, the Hamiltonian operator is set up for the system, accounting for the kinetic and potential energy of the particles constituting the system, then inserted into the Schrödinger equation. The resulting partial differential equation is solved for the wave function, which contains information about the system.** |
| **Time-independent equation** |
| **(31) The time-dependent Schrödinger equation described above predicts that wave functions can form standing waves, called stationary states (also called "orbitals", as in atomic orbitals or molecular orbitals). These states are important in their own right, and if the stationary states are classified and understood, then it becomes easier to solve the time-dependent Schrödinger equation for any state. Stationary states can also be described by a simpler form of the Schrödinger equation, the time-independent Schrödinger equation. (This is only used when the Hamiltonian itself is not dependent on time explicitly. However, even in this case the total wave function still has a time dependency.)** |
| **Time-independent Schrödinger equation (general)** |
| **(32) E Ψ = H ^ ⊠ Ψ {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi } {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi }** |
| **In words, the equation states:** |
| **(33) When the Hamiltonian operator acts on a certain wave function Ψ, and the result is proportional to the same wave function Ψ, then Ψ is a stationary state, and the proportionality constant, E, is the energy of the state Ψ.** |

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| **(34) The time-independent Schrödinger equation is discussed further below. In linear algebra terminology, this equation is an eigenvalue equation.** |
| **(35) As before, the most famous manifestation is the non-relativistic Schrödinger equation for a single particle moving in an electric field (but not a magnetic field):** |
| **Time-independent Schrödinger equation (single non-relativistic particle)** |
| **(36) E Ψ ( r ) = [ − ℏ 2 2 μ ∇ 2 + V ( r ) ] Ψ ( r ) {\displaystyle E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )} E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )** |
| **Gravitational waves: (Wikipedia)** |
| **(37) In theory, the loss of energy through gravitational radiation could eventually drop the Earth into the Sun. However, the total energy of the Earth orbiting the Sun (kinetic energy + gravitational potential energy) is about 1.14×1036 joules of which only 200 Watts (joules per second) is lost through gravitational radiation, leading to a decay in the orbit by about 1×10−15 meters per day or roughly the diameter of a proton. At this rate, it would take the Earth approximately 1×1013 times more than the current age of the Universe to spiral onto the Sun. This estimate overlooks the decrease in r over time, but the majority of the time the bodies are far apart and only radiating slowly, so the difference is unimportant in this example.** |
| **More generally, the rate of orbital decay can be approximated by[39]** |
| **(38) d r d t = − 64 5 G 3 c 5 ( m 1 m 2 ) ( m 1 + m 2 ) r 3 , {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,} {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,}** |
| **where r is the separation between the bodies, t time, G Newton's constant, c the speed of light, and m1 and m2 the masses of the bodies. This leads to an expected time to merger of [39]** |
| **t = 5 256 c 5 G 3 r 4 ( m 1 m 2 ) ( m 1 + m 2 ) . {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.} {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.}** |
| **(39) For example, a pair of solar mass neutron stars in a circular orbit at a separation of 1.89×108 m (189,000 km) has an orbital period of 1,000 seconds, and an expected lifetime of 1.30×1013 seconds or about 414,000 years. Such a system could be observed by LISA if it were not too far away. A far greater number of white dwarf binaries exist with orbital periods in this range. White dwarf binaries have masses in the order of the Sun, and** |

**diameters in the order of the Earth. They cannot get much closer together than 10,000 km before they will merge and explode in a supernova which would also end the emission of gravitational waves. Until then, their gravitational radiation would be comparable to that of a neutron star binary.**

**(40) By substituting value of dr/dt in Schrödinger's wave equation we get the quantum wave equation for gravitational waves.**

In order to have (H3) and (H4) hold, $m, \mu_0, \mathrm{K}, \mathrm{b}, \delta$, and $\omega$ need to satisfy

$$\frac{\omega_1}{(m+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1+b_1}+\sigma)^2}+\frac{2\omega_1}{([m_1]_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1+b_1}+\sigma_1)^3}<1$$

And

$$\frac{1}{(m_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1-b_1})^2}+\frac{2\omega_1}{(m_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1-b_1})^3}>1$$

When the inequalities (33) and (34) hold there is a unique equilibrium $\mathrm{Q}^*$ which can be solved from the following equation:

$$(m_1+[\mu_0]_1+1+\tfrac{K_1 b_1}{\delta_1-b_1}+(m_1+[\mu_0]_1+\tfrac{K_1 b_1}{\delta_1-b_1}+\sigma_1)\,\mathrm{Q}_1^*)^3$$

$$=m_1+[\mu_0]_1+3+\tfrac{K_1 b_1}{\delta_1-b_1}(1+\mathrm{Q}^*)(1+\omega_1\mathrm{Q}_1^*$$

Now

$$\left[\mathrm{a}_1^2\right]-4\left[\mathrm{a}_2\right]\mathrm{e}=\left(\frac{2K_1 b_1}{\delta_1-b_1}\right)^2+\frac{4K_1 b_1\sigma_1\,\mathrm{Q}_1^*}{(\delta_1-b_1)(1+\mathrm{Q}_1^*)}+\frac{\sigma_1^2\,\mathrm{Q}_1^{*2}}{(1+\mathrm{Q}_1^*)^2}$$

$$\frac{4\,b_1\sigma_1\,(\delta_1-\mathrm{b})\mathrm{Q}^{*}}{\delta_1\,(1+\mathrm{Q}_1^{*})^{2}}-\frac{4\,(\delta_1-b_1)^{2}(1-\omega_1)\mathrm{Q}_1^{*}}{\delta_1\,(1+\omega_1\mathrm{Q}_1^{*})(1+\mathrm{Q}_1^{*})}\mathrm{P}^{*}(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_1\,ae^{mu}_1\;\;[\mu_0]_1-\frac{a_1}{a_1+1}$$

Hence if $[\mathrm{a}_1^2>4(a_2)_1\mathrm{e}$ and

$$[a_1]=\frac{2K_1b_1}{\delta_1-b_1}+\frac{\sigma_1\,\mathrm{Q}_1^{*}}{1+\mathrm{Q}_1^{*}}>2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_1]_1=[f_{1*}]_1+[f_{2*}]_1+[f_{1*}]\,\mathrm{P}_1^{*}$$

$$[a_2]_1=\left[\mathrm{f}_{2*}^{1}\right]\mathrm{P}_1^{*}-\frac{\mathrm{h}_*^1}{\mathrm{h}_*}\mathrm{P}^{*}(0)$$

$$\mathrm{e}=[\,\mathrm{g}_*^1]_1\,\mathrm{Q}_1^{*}$$

$$\eta_1(\mathrm{t})=\left[\mathrm{h}_*\right]_1\beta_1\left(\theta_1(\mathrm{t})\right)-\mu_1(\theta(\mathrm{t})),$$

The system(31)can be written as the following system:

$$[\mathrm{M}(\mathrm{t})]\left[\mathrm{a}_1\right]M_1-[a_2]_1W_1-[\mathrm{n}(\mathrm{t})]M_1$$

$$w_1(t) = eM_1$$

Write

$$A=\begin{pmatrix} -[a_1]_1 & -[a_2]_1 \\ e & 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_1 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_1 \\ W_1 \end{pmatrix}$$

The solution of (32)is

$$X_1(t) = e^{A,t}[X]_1(0) + \int_0 e^{At}$$

## Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

## (Module Three)

**Abstract: Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-435) 1990. We demonstrate its applicability cosmology/ quantum computation economics et al. As stated in the foregoing we shall apply model to various cosmological problem and show the relevance to the same. . To exemplify model can be applied to the following situational significations: Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new**

**source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 Masaru Shibata and Takashi Nakamura** perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. **Authors** use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. 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Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^{\infty} \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(g_3)Q_3$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_3(0) = 0, \quad g_3(P_3) > 0, \quad P_3 \; 0, \qquad \lim_{P_3 \to \infty} g_3(P_3) = (K_3)_3 \; \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_3 + (\rho_a)_3 + D_3(a_3, P_3, Q_3)\rho_3 + 0, \qquad 0 < a_3 \; \infty,$$

$$\rho_3(t, 0) = \int_0^{\infty} n_3(a_3, P_3, Q_3)\rho_3(t_3, a_3)da_3$$

Where $D_3$ is the mortality function and $n_3$ is the fertility function.

Put $_3$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_3}{dt}\right]_3 = x_3 \left[\frac{dQ_2}{dt}\right]_3 - x_2 \left[\frac{dQ_3}{dt}\right]_3$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3 \left[\frac{dQ_3}{dt}\right]_3 - x_4 \left[\frac{dQ_2}{dt}\right]_3$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5 \left[\frac{dQ_2}{dt}\right]_3 - x_6 \left[\frac{dQ_3}{dt}\right]_3$$

Where $x_3, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_3$ and $g_3(P_3)Q_3$

Now when you add the term $(b_3Q_3)$ with $g_3(P_3)(Q_3)$ we are infect see is that $(b_3Q_3)$ is credited incrementally part of $g_3(P_3)(Q_3)$ or instant nearly, when subtracted from $\frac{dQ_3}{dt}$ , it means that $g_3(P_3)(Q_3)$ is slowly eaten by $\frac{dQ_3}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_3 = \mu_3(a_3) + (f_3)_3(P_3)_3 + (f_2)_3(Q)_3 \quad (2)$$

$$n_3 = \beta_3(a_3)h_3(Q_3), \quad (3)$$

Where $\mu_3$the age- dependent debit rate whit is $\mu_3(a_3)$ $\forall a > 0$; $(f_3)_3$ is the density- dependent death rate with $f_3(0) = 0$, $(f_3^3)_3(P_3)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_3)_3(P_3) = (k_3)_3$ ; $f_2$ is the functional response of creditors to prey with $[f_2]_3(0) = 0, (f_2^3)_3(Q) > 0, \forall Q$ , and $\lim_{Q\to\infty}(f_2)_3(Q_3) = [k_2]$ ; $\beta$ is the age dependent birth rate with $\beta(a)$ $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 3, h^3(Q)$ $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ .

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_3 + (\rho_a)_3 = -\mu_3(a_3) + (f_3)_3P_3) + \big((f_2)_3(Q_3)\big)(\rho_3)_3 \quad (4)$$

$$\rho_3(t_3, 0) = h_3(Q_3)\int_0^\infty \beta_3(a_3)\rho_3(t, a_3)da, \quad (5)$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(P_3)Q_3 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_3$ and $Q_3$ of (4-6) satisfy

$$0 \le \rho_3(t, a_3) \le \rho_3^*(t, a_3),$$

$$0 \le Q_3(t) \le Q_3^*(t),$$

Where $\rho_3^*(t, a_3)$ and $Q_3^*(t)$, are solutions of

$$(\rho_t^*)_3 + (\rho_a^*)_3 = -(\mu_3(a_3) + (f_3)_3(P_3^*) + (f_2)_3 Q_3^*))\,\rho_3^*, \quad (7)$$

$$\rho_3^*(t, 0) = \int_0^\infty \beta_3(a_3)\,\rho_3^*(t, a_3)da_3 \quad (8)$$

$$\frac{dQ_3^*}{dt} = -b_3 Q_3^* + g_3(P_3^*)Q_3^*, \quad (9)$$

With $P_3^*(t) = \int_0^\infty \rho_3^*(t, a_3)da_3$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_3^*(t)}{dt} = P_3^*(t)\big(r_3 - (f_3)(P_3^*) - (f_2)_3(Q^*)\big), \quad (10)$$

$$\frac{dQ_3^*(t)}{dt} = Q_3^*(t)(-b_3 + g_3(P_3^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_3(a)e^{-\int_0^a \mu_3(\tau)dt}e^{-ra}\,da_3 = 1. \qquad (12)$$

If $r < 0$, the trivial solution $P_3^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(t)dt}\,da_3 > 1, \qquad (H1)$$

Which guarantees $r_3 > 0$.

*Remark.* If $\beta_3(a) = [\beta_0]_3 a_3 e^{-\alpha_3 a_3}, \mu_3(a_3) \equiv \mu_3$ then the solution of (12) is

$$r_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_3]_3(P_3) = 0, [f_2](Q_3) = \lambda_3 Q_3\ h_3(Q_3) = 1, \qquad and\ \ g_3(P_3) = c_3 P_3$$

$$[\rho_3]_3 + [\rho_a]_3 = -(\mu_3 + \lambda_3 Q_3)\rho_3$$

$$\rho_3(t,0) = \int_0^\infty [\beta_0]_3 a_3 e^{-\alpha_3 a_3}\rho_3(t,a_3)\,da_3$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + c_3 P_3 Q_3$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_3}{dt} = \omega_3 P_3 - \lambda_3 P_3 Q_3$$

$$\frac{dQ_3}{dt} = b_3 Q_3 - c_3 P_3 Q_3$$

Where $\omega_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$

## THE EQUILIBRIUM $(\rho_0)_3(a_3), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_3(a_3), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_3 + (\rho_a)_3 = -\big(\mu_3(a_3) + (f_3)_3(P_3)\big)\rho,$$

$$\rho_3(t,0) = \int_0^\infty \beta_3(a_3)\rho_3(t,a_3)da_3 \qquad (13)$$

The solution of (13) has the form of

$$\rho_3(t,a_3) = Ce_3^{r(t-a_3)} - \int_0^a (\tau)d\tau - \int_0^3 (f_3)_3\big(P_3(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_3(t) = C_3 e^{r_3 t - \int_0^t (f_3)_3(P_3(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_3^{-3}(r)$ if $k_3 > r$. This is equivalent to

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{(k_3)_3 a_3} da_3 < 1.$$

In fact, if we define

$$R_3(s_3) = \int_0^{\infty} \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau)d\tau} e^{-s_3 a_3}\, da_3$$

$R(s)$ is a decreasing function. $R_3(r_3) = 1$ And $R_3(k_3)_3 < 1$ imply $r_3 < (k_3)_3$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_3(a_3), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{k_3 a_3} da_3$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_3 \equiv \int_0^{\infty}(\rho_0)_{33}(a_3)da_3$ *satisfies*

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{-(f_3)_3(P_0)a_3} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_3(a_3), 0)$

Let $u_3(t, a_3) = \rho(t, a_3) - \rho_0(a_3)$ Then

$$N_3(t) \equiv \int_0^{\infty} u_3\,(t, a_3)da_3 = \int_0^{\infty} \rho_3(t, a_3)da_3 - \int_0^{\infty} (\rho_0)_3(a_3)da = P_3(t) - P_0)_3$$

And

$$(u_3)_3 + (u_a)_3 = (\rho_t)_3 + (\rho_a)_3 - (\rho_0^3)_3(a_3)$$

$$= -\mu_3(a_3) + (f_3)_3(P_3) + (f_2)_3(Q_3))\,(u_3 + (\rho_0)_3) + \mu_3(a_3)(\rho_0)_3 + (f_3)_3(P_0)_3(\rho_0)_3$$

$$= -\mu_3(a_3)u_3 - (f_3)_3(P_3) + (f_2)_3(Q_3))\,u_3 - (\rho_0)_3\,(f_3)_3(P_3) - (f_3)_3(P_0)_3$$

$$+ (f_2)_3(Q_3)), \qquad (15)$$

Where $(\rho_0^3)_3(a_3)$ denotes $d(\rho_0)_3(a_3)/da_3$

Expansion $(f_3)_3(P_3)$ and $(f_2)_3(Q_3)$ at $(P_0)_3$ and 0 yield respectively

$$(f_3)_3(P_3) = (f_3)_3((P_0)_3 + N_3)$$

$$= (f_3)_3(P_0)_3 + (f_3^3)_3(P_0)_3N_3 + o_3(N_3), \qquad (16)$$

$$(f_2)_3(Q_3)$$

$$= (f_2^3)_3(0)Q + o_3(Q_3) \qquad (17)$$

Write $f_3(P_0) = f_{30}, f_3'(P_0) = f_{30}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_3)_3 + (u_a)_3 \approx -(\mu_3(a_3) + (f_{30})_3)u - (f_{30}^3)_3(\rho_0)_3(N_3) - (f_{20}^3)_3(\rho_0)_3(Q_3)$$

$$u_3(t,0) = \rho_3(t,0)(\rho_0)_3(0)$$

$$= h_3(Q_3)\int_0^\infty \beta_3(a_3)(u_3(t,a) + (\rho_0)_3(a_3)da_3 - (\rho_0)_3\,(0)$$

$$= h_3(Q_3)\int_0^\infty \beta_3(a_3)(u_3(t,a_3)da_3 + (h_3(Q_3) - 3)(\rho_0)_3\,(0) \qquad (18)$$

Expanding

$$h_3(Q_3) = 1 + (h^3)_3(0)Q_3 + 0(Q_3) \equiv 1 + h_0'Q_3 + 0Q_3$$

And substituting it into (38), we arrive at

$$u_3(t,0) \approx \int_0^\infty \beta_3\,(a_3)u_3(t,a_3)da_3 + (h_{a_3}')(\rho_0)_3Q_3$$

The linearization is the following :

$$[u_3] + [u_3] = -(\mu_3(a_3) + [t_{30}]_3 u_3 - [\alpha_3]_3(a_3)N_3 - [\alpha_2]_3(a_3)Q_3$$

$$u_3(t,0) = \int_0^\infty \beta_3(a_3)u_3(t,a_3)da_3 + [\xi_3]Q_3 \qquad (19)$$

$$\frac{dQ_3}{dt} = -b_3 Q_3$$

where

$$[\alpha_3]_3(a_3) = [t_3^3 a]_3[p,o]_3(a_3)d2(a_3) = [f_{20}^3]_3(\rho\, o\,)_3(a_3)\xi = [\mathrm{h\ o^3}]_3[\rho\mathrm{o}]_3(\mathrm{o}) \text{ and } \mathrm{b}_3 = \mathrm{b}_3 - g_3(P_3)$$

If $g_3(P_0) > b$, clearly, $(\rho_0)_3(a_3), 0)$ is unstable and the predator population is persistent. Assume $g_3(P_0)_3 < b$.

Then, $Q_3(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_3 + u_a = -(\mu_3(a_3) + (f_{30})_3)u_3 - (\alpha_3)_3(a_3)N_3$$

$$= -(\mu_3(a_3) + f_{30})u_3 - \alpha_3(a_3)\int_0^\infty u_3(t,a_3)da_3$$

$$u_3(t,0) = \int_0^\infty \beta_3\,(a_3)u_3(t,a_3)da_3$$

Using the method developed by Cushing [2, 3], we suppose

$$u_3(t,a_3) = \tilde{\mathrm{u}}_3(a_3)^{e^{\gamma(t-a_3)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_3}\, da_3$. Then,

A solution of the first differential equation for $\tilde{u}(a_3)$ gives,

$$\tilde{u}_3(a_3) = e^{-\int_0^a \mu_3(\tau_3)d\tau - (f_{30})_3 a}\left(\tilde{u}_3(0) - w\int_0^a e^{\int_0^\tau \mu_3(s_3)ds + (f_{30})}\tau_3\,(\alpha_3)_3(\tau_3)e^{\gamma_3\tau}d\tau\right).$$

Substituting the expression (13) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\right)\tilde{u}_3(0)$$

$$+w\int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\int_0^a (\alpha_3)_3(\tau_3)e^{-\gamma_3\tau_3}e^{(f_{30})_3\tau_3 + \int_0^\tau \mu_3(s_3)ds}d\tau_3\, da_3$$

$$= 0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}da\right)(\tilde{u}_3(0))_3$$

$$+\left(-3 - \int_0^\infty e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\int_0^a e^{(f_{30})_3\tau_3 + \int_0^\tau \mu_3(s_3)ds + y\tau}\alpha_3(\tau)d\tau\, da\right)w = 0.$$

(23)

Define $Y(\gamma, a) = e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(t_3)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_3\,(a_3)Y_3(\gamma_3,a_3)da_3\right)u_3(0)$$

$$+w\int_0^{\infty}\beta_3(a_3)Y_3(\gamma_3,a_3)\int_0^{a}(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\,da_3=0,$$

(24)

$$\left(\int_0^{\infty}Y_3(\gamma_3,a_3)da_3\right)\tilde{u}_3(0)$$

$$+\left(-1-\int_0^{\infty}Y_3(\gamma_3,a_3)\int_0^{a}(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\,da_3\right)w=0.$$

Equqtion (20) has a solution of the form of $u_3(t,a_3)=u_3(a_3)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_3(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_3(\gamma_3,a_3)=\int_0^{a}(\rho_0)_3(\tau_3)Y^{-3}\,(\gamma_3,a_3)d\tau_3$. We obtain the characteristic equation

$$f_{30_3}^{l}\left(\int_0^{\infty}Y_3(\gamma_3,a_3)da_3\right)\left(\int_0^{\infty}\beta_3(a_3)Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$+\left(1-\int_0^{\infty}\beta_3(a_3)Y_3(\gamma_3,a_3)da\right)\left(1+(f_{30}^{l})_3\int_0^{\infty}Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_3(a_3)=(\rho_0)_3(0)e^{-\int_0^{a}\mu_3(\tau_3)d\tau_3-(f_{30})_3a_3}$$

$$I_3(\gamma_3, a_3) = \int_0^a (\rho_0)_3(\tau_3) e^{\gamma_3\tau_3 + (f_{30})_3\tau + \int_0^\tau \mu_3(s_3)ds_3} d\tau$$

$$= \frac{(\rho_0)_3(0)}{\gamma_3}(e^{\gamma_3 a_3} - 1)$$

Let

$$E_3(a_3) = e^{-(f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau_3}$$

Then $[Y_3]_3[\gamma_3]_3, a_3) = E_3(a_3)e^{-\gamma_3 a_3}$. by a simple calculation,

$$\int_0^\infty \beta_3(a_3)Y_3(\gamma_3, a_3)I(\gamma_3, a_3)da = \frac{[\rho_0]_3(0)}{\gamma}\left(1 - \int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}\ da\right);$$

$$\int_0^\infty Y_3(\gamma_3, a_3)I(\gamma_3, a_3)da = \frac{[\rho_0]_3(0)}{\gamma_3}\left(\int_0^\infty E_3(a_3)e^{-\gamma_3 a_3}\ da_3\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_3(0)}{\gamma_3}[f_{30}^3]\int_0^\infty E_3(a_3)e^{-\gamma_3 a_3}da_3\ (\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3)$$

$$+\left(1 - \int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right).$$

$$\left(1 + \frac{(\rho_0)_3(0)}{\gamma_3}(f_{30}^{'})_3\left(\int_0^\infty E_3(a_3)da_3 - \int_0^\infty E_3(a_3)e^{-\gamma_3 a_3}da_3\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right)(1+\frac{(f_{30}^{'})_3}{\gamma_3}(\rho_0)_3(0)\int_0^{\infty}E_3(a_3)da_3)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_3(a_3)Y_3(\gamma_3 a_3)da_3=1-\int_0^{\infty}\beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\neq 0,$$

Since $\int_0^{\infty}\beta_3(a_3)Y_3(\gamma_3 a_3)\int_0^{a}\alpha_3(\tau_3)Y^{-3}(\gamma_3,\tau_3)d\tau_3 da_3>0$. Therefore, the characteristic solution is

$$\gamma_3=-(f_{30}^{'})_3(\rho_0)_3\int_0^{\infty}E_3(a_3)da_3.$$

Since $f_3'(P_0)_3>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_3),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_3),0)$ which is related to the extinction of debit summation types

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a),Q^*)$. It must satisfy the equations:

$$[p_a^*]_3 + (\mu_3(a_3) + [f_3]_3(P_3^*) + [f_2]_3(Q_3^*))P_3^* = 0, \quad (26)$$

$$p_3^*(0) = h_3(Q_3^*)\int_0^\infty \beta_3(a_3)p_3^*(a_J)p_3^*(a_3)da_3 \quad (27)$$

$$g_3(P_3^*) - b_3 = 0, \quad (28)$$

Where $P_3^*=\int_0^\infty p_3^*(a_3)\,da_3$

Equation (28) yields $P^* = g^{-3}(b)$,

From (26),

$$p_3^*(a) = p_3^*(0)e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}$$

Substituting into (27), we need to solve

$$\left(1 - h_3(Q_3^*)\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}da_3\right)p_3^*(0) = 0$$

For $Q_3^*$

Fix $p_3^*(0)$ and $P_3^*$ Let

$$H_3(z_3) = 1 - h_3(z_3)\int_0^\infty \beta_3(a_3)\,e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3}da_{,3}$$

And assume that

$$[k_3]_3\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(g_3^{-3}(b_3))a_3 - [k_2]_3a_3}\,da_3 < 1. \quad (H3)$$

Then

$$\lim_{z\to\infty}[H_3(z_3)] > 0.$$

Clearly,

$$\frac{dH_3}{dz_3} = h_3(z)([f_2^{'}]_3(z_3)) \int_0^{\infty} a_3\beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - f_3(P_3^*)a_3 - f_2(z_3)a_3}\, da_3$$

$-h^{'}(z_3)\int_0^{\infty}\beta_3(a_3)\, e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3}\quad da_3$

By the assumptions on $\left[f_2^3\right]_3 and\ [h^3]^3 dH_3/dz > 0$. Hence, if

$$H_3(0) = 1 - \int_0^{\infty}\beta_3(a_3)e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3}da_3 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_3(0) < 0$ follows if $g_3^{-}(b_3) < [f_3^{-3}]_3(r_3)$, where $r_3$ is the solution of (12).

Indeed, since

$$R_3(s_3) = \int_0^{\infty}\beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 e^{-y\frac{a}{3}}}da_3$$

Is a decreasing function of $s_3$ when $P_3^* = [g^{-3}]_3(b_3) < [f_3^{-3}]_3(r_3)$, that is $[f_3]_3(P_3^*) < r$, $R\left(f_3(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_3(g^{-3}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3\left(g_3^{-3}(b)\right)a_3} > 1. \quad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_3^*(a_3), Q_3^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_3^*(a_3), Q_3^*)$ is more complicated. We again consider perturbations $m_3(t_3 a_3) = p_3(t_3, a_3) - p_3^*(a_3)$ and $W_3(t) = Q_3(t) - Q^*$.Then

$$M_3(t) \equiv \int_0^\infty m_3(t, a_3) da_3 = P_3(t) - P_3^*,$$

$$[m_3]_3 + [m_a]_3 = [p_3]_3 + [p_a]_3 - \left[p^*\right]_3(a_3) \approx -(\mu_3(a_3) + [f_{3*}]_3 + [f_{2*}]_3)m_3$$

$$-[f_{3*}^3]_3 p_3^*(a_3)M - [f_{2*}^3]_3 p_3^*(\mathrm{a})W_3,$$

$$m_3(t,0)p_3(t,0) - p_3^*(0) \approx h_* \int_0^\infty \beta_3(a_3) m_3(t, a_3)\, da + \frac{[h_*^3]_3}{[h_*]_3}\, p_3^*(0) W_3$$

$$\frac{dW_3}{dt} = \frac{dQ_3}{dt} \approx [g_*^3]_3 Q_3^* M_{,3}$$

Where $[f_{3*}]_3$ denoting $[f_3(P^*)]_3$, $[f_{3*}]_3$ denoting $[f_2(P^*)]_3$, $\left[f_{3*}^{'}\right]_3$ denoting $\left[f^{'}{}_3(P^*)\right]_3 \left[f_{2*}^{'}\right]_3$ denoting $[f_2]_3$ $(Q_3^*,$ $h_3^*$ denoting $h_*^{'}$ denoting $h_3^{'}$ $(Q_3^*)$, and $\left[g_*^{'}\right]_3$

Denoting $\left[g^{'}\right]_3(P_3^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_3]_3 + [m_a]_3 = -(\mu_3(a_3) + [f_{3*}]_3 + [f_{2*}]_3)m_3 - [f'_{3*}]_3[p_3^*](a_3)M_3 - [f'_{2*}]_3 p_3^*(a_3)W_3$$

$$m_3(t,0) = [h_*]_3 \int_0^\infty \beta_3(a_3)m_3(t,a_3)da_3 + \frac{[h_*^3]_3}{[h^*]_3}[p^*]_3(0)\,W_{,3} \qquad (30)$$

$$\frac{dW_3}{dt} = [g_*^3]_3[Q_3^*]M_3$$

Integrating $[m_3]_3(t,a_3)$ with respect to a from 0 to ∞, we have

$$\frac{dM_3(t)}{dt} = \int_0^\infty [mt]_3(t,a_3)da_3$$

$$= -\int_0^\infty [m_a]_3\,(t,a_3)da_3 - \int_0^\infty \mu(a_3)m(t,a_3)da_3$$

$$-\left([f_{3*}]_3 + [f_{2*}]_3 + [f'_{3*}]_3\int_0^\infty p_3^*(a_3)da_3\right)M_3 - [f_{2*}^3]_3\int_0^\infty p_3^*(a_3)da_3 W_3$$

$$= m_3(t,0) - \int_0^\infty \mu_3(a_3)\,m_3(t,a_3)da_3$$

$$-\left([f_{3*}]_3 + [f_{2*}]_3 + [f'_{3*}]_3 P_3^*\right)M_3 - [f'_{2*}]_3 P_3^* W_{,3}$$

From (30), we arrive at the equations

$$\frac{dW_3(t)}{dt} = \int_0^\infty ([h_*]_3\beta_3(a_3) - \mu_3(a_3))m_3(t,a_3)da_3$$

$$-\left([f_{3*}] + [f_{2*}]_3 + [f'_{3*}]_3 P_3^*\right)M_3 + \left(\frac{[h'_*]_3}{[h_*]^3}\,P_3^*(0) - f'_{2*}P_3^*\right)W_3,$$

(33)

$$\frac{dW_3(t)}{dt}\left[g^{'}{}_{*}\right]_3 Q_3^{*}\, M_{.3}$$

By the mean value theorem,

$$\int_0^{\infty}([h_{*}]_3\beta_3(a_3)-\mu_3(a_3))m_3(t,a_3)da_3=([h_{*}]_3\beta_3(\theta_3)-\mu_3(\theta_3))M_{,3}$$

Where $0<\ \theta_3=\theta_3(t)<\infty$.

And if $a_3^2\,[a_3^2]_3<4\,[a_2]_3<4\,a_2 e$,

Re $\lambda_3\left[\lambda_3\right]_3=\mathrm{Re}\lambda_2\left[\lambda_2\right]_3=-\frac{3}{2}a_3[a2\ ]_3<0.$

Let $\Delta=\max_{2<u<\infty}\left|h_{*}\,[h_{*}]_3\beta_3(a_3)-\mu_3(a_3)\right|$. Then$|n_3(t)|$ 2Δ∀t 0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium($\rho^{*}(a)$, $\boldsymbol{Q}^{*}$) is asymptotically stable if either $a_3^2<4\,a_2 e$ and $[a_3]>2\Delta$,or $[a_3^2]$ 4$[a_2]$e and $(a_3-\sqrt{(a_3^2)_3-4(a_2)_3 e})>2\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a)=\mu_0+\left(\frac{a}{a+3}\right)$ and $\beta_3(a_3)=ae^{-ma}$ Then $\Delta=[\mu_0]+1$ Simple calculation leads to

$$\int_0^{\infty} \beta_3(a_3)\mathrm{e}^{-\int_0^{\mathrm{a}} \mu_3(\tau_3)+d\tau_3 \mathrm{da}} = \frac{3}{(\mathrm{m}+\mu_0+1)^2}+\frac{2}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_3 + [\mu_0]$ 0.5. Assume$[f_3](x_3)= K_3x_3$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_3(x_3)= [0x_3]/(3 + x_3)$; the functional responses $[[f_2](x) = [0]\, x_3/(1 + x_3)$and $h_3(x_3)$=$(1+\omega_3 x_3)/(1 + x_3)$(See Holling [7] and Rashevsky [13])

Equation (34),now has the form of

$$\frac{m_3 + [\mu_0] + [\mathrm{K}][P_0])}{[m_3]_3 + (\mu_0)_3 + [\mathrm{K}][P_0])_3^3\ [\mathrm{K}]_3[P_0]_3^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(3 + P_0) < b_3$i e $[P_0] < b/(\delta_3 - b_3)$then the equilibrium$(P_0)$ $[P_0]_3(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_3$

$[f_{3*}] = K_3\ \mathrm{P}^*\ [f_{2*}] = \frac{\sigma \mathrm{Q}_3^*}{3+\mathrm{Q}_3^*}$ $[h_*] = \frac{3+\omega_3 \mathrm{Q}_3^*}{3+\mathrm{Q}_3^*}$

$[f_{3*}] = K_3\ [\mathrm{f}_{2*}^3] = \frac{\sigma_3}{(3+\mathrm{Q}_3^*)^2}[\ \mathrm{g}_*^3] = \frac{\sigma_3}{(3+\mathrm{P}_3^*)^2}$ $[\mathrm{h}_*^3] = \frac{\omega_3-3}{(3+\mathrm{Q}_3^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_3}{(m+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma)^2}+\frac{2\omega_3}{([m_3]_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma_3)^3}<1$$

And

$$\frac{1}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^2}+\frac{2\omega_3}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}+(m_3+[\mu_0]_3+\frac{K_3b_3}{\delta_3-b_3}+\sigma_3)\,Q_3^*)^3$$

$$=m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}(1+Q^*)(1+\omega_3 Q_3^*$$

Now

$$\left[a_3^2\right]-4\left[a_2\right]e=\left(\frac{2K_3b_3}{\delta_3-b_3}\right)^2+\frac{4K_3b_3\sigma_3\,Q_3^*}{(\delta_3-b_3)(1+Q_3^*)}+\frac{\sigma_3^2\,Q_3^{*2}}{(1+Q_3^*)^2}$$

$$\frac{4\,b_3\sigma_3\,(\delta_3-b)Q^*}{\delta_3\,(1+Q_3^*)^2}-\frac{4\,(\delta_3-b_3)^2(1-\omega_3)Q_3^*}{\delta_3\,(1+\omega_3Q_3^*)(1+Q_3^*)}P^*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_3\,ae_3^{mu}\;\;[\mu_0]_3-\frac{a_3}{a_3+3}$$

Hence if $[a_3^2>4(a_2)_3e$ and

$$[a_3]=\frac{2K_3b_3}{\delta_3-b_3}+\frac{\sigma_3\,Q_3^*}{1+Q_3^*}>2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_3]_3 = [f_{3*}]_3 + [f_{2*}]_3 + [f_{3*}]P^*_3$$

$$[a_2]_3 = \left[f^3_{2*}\right]P^*_3 - \frac{h^3_*}{h_*}P^*(0)$$

$$e=[g^3_*]_3 Q^*_3$$

$$\eta_3(t) = \left[h_*\right]_3 \beta_3\left(\theta_3(t)\right) - \mu_3(\theta(t)),$$

The system(33)can be written as the following system:

$$[M(t)]\left[a_3\right]M_3 - [a_2]_3 W_3 - [n(t)]M_3$$

$$w_3(t) = eM_3$$

Write

$$A=\begin{pmatrix} -[a_3]_3 - [a_2]_3 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_3 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_3 \\ W_3 \end{pmatrix}$$

The solution of (32)is

$$X_3(t) = e^{A,t}\,[X]_3(0) + \int_0 e^{At}$$

**This equation is useful in following type of problems . For actual model kindly see various papersThere is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant Λ, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive Λ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the Λ>0 generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible. Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of**

**(e) GW signals from BBHs with component spins χ1,2∈[−0.99,0.99] aligned (e&eb) with the orbital angular momentum, component masses m1,2∈[2,48]M⊙, and (e&eb) total mass Mtotal≤50 M⊙. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins >+0.6 and component masses ≳20 M⊙. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201 STOP Advanced LIGO may be the first experiment to (e) detect gravitational waves. Through superradiance of stellar black holes, it may also be (=) the first experiment to discover the QCD axion with (e&eb) decay constant above the GUT scale. When an axion's Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e) the black hole, forming a "gravitational atom." Through the superradiance process, the number of axions occupying (e) the bound levels grows exponentially, extracting (e) energy and angular momentum from the black hole. Axions transitioning between levels of the gravitational atom and (e&eb) axions annihilating to (e) gravitons can produce (eb) observable gravitational wave signals. The signals are long-lasting, monochromatic, and can be distinguished from (e) ordinary astrophysical sources. Authors estimate up to O (1) transition events at aLIGO for (e) an axion between 10^-11 and (e&eb) 10^-10 eV and up to 10^4 annihilation events for (e) an**

**axion between 10^-13 and 10^-11 eV. In the event of a null search, aLIGO can constrain (e) the axion mass for (e) a range of rapidly spinning black hole formation rates. Axion annihilations are also promising for (e) much lighter masses at future lower-frequency gravitational wave observatories; the rates have (e) large uncertainties, dominated by (e) supermassive black hole spin distributions. Projections for aLIGO are (=) robust against perturbations from (e) the black hole environment and account for (e) updated exclusion on (e&eb) the QCD axion of 6*10^-13 eV < ma < 2*10^-11 eV suggested by (e) stellar black hole spin measurements. Subjects: High Energy Physics - Phenomenology (hep-ph); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 91, 084011 (2015) DOI: 10.1103/PhysRevD.91.084011Cite as:arXiv:1411.2263 [hep-ph] (or arXiv:1411.2263v3 [hep-ph] for this version) Discovering the QCD Axion with Black Holes and Gravitational Waves Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang STOP Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by**

**(e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave signal. The signal sweeps upwards in frequency from 35 to 250 Hz with (e&eb) a peak gravitational-wave strain of (e) 1.0×10−21. It matches the waveform predicted by general relativity for (e) the inspiral and merger of a pair of black holes and the ringdown of the resulting (eb) single black hole. The signal was observed with a matched-filter signal-to-noise ratio of (e) 24 and a false alarm rate estimated to be (=) less**

**than 1 event per 203 000 years, equivalent to (=) a significance greater than 5.1σ. The source lies at (eb) a luminosity distance of 410+160−180 Mpc corresponding to (e&eb) a redshift z=0.09+0.03−0.04. In the source frame, the initial black hole masses are (=) 36+5−4M⊙ and 29+4−4M⊙, and the final black hole mass is (=) 62+4−4M⊙, with (e&eb) 3.0+0.5−0.5M⊙c2 radiated in (eb) gravitational waves. All uncertainties define (eb) 90% credible intervals. These observations demonstrate (eb) the existence of binary stellar-mass black hole systems. Observation of Gravitational Waves from a Binary Black Hole Merger B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 061102 – Published 11 February 2016 This is the first direct detection of gravitational waves and (e&eb) the first observation of a binary black hole merger. Received 21 January 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.061102 STOP This is the second in a series of three papers on (e&eb) systematic analysis of rank 1 Coulomb branch geometries of (e) four dimensional N=2 SCFTs. In the first paper authors developed a strategy for (e) classifying physical rank-1 CB geometries of N=2 SCFTs. They show how to carry out this strategy computationally to construct (eb) the Seiberg-Witten curves and (e&eb) one-forms for all the rank-1 SCFTs. Explicit expressions are given for all cases, with the exception of the Nf=4 SU (2) gauge theory and (e&eb) the En SCFTs which were previously constructed. Classification includes all known rank-1 theories plus (+) a new one with (e&eb) an abelian flavor group, plus (+) nine additional theories whose existence is more speculative. Four of those, reported in first paper, depend on (=) the assumption of new frozen rank-1 SCFTs. Here authors also show that the assumption of the existence of certain rank-0 N=2 SCFTs leads to (eb) five additional consistent rank-1 CB geometries. Subjects: High Energy Physics - Theory (hep-th) Cite as:arXiv: 1601.00011 [hep-th] (or arXiv: 1601.00011v2 [hep-th] for this version Geometric constraints on the space of N=2 SCFTs II: Construction of special Kähler geometries and RG flows Philip C. Argyres, Matteo Lotito, Yongchao Lü, Mario Martone (Submitted on 31 Dec 2015 (v1), last revised 6 Jan 2016 (this version, v2)) STOP Authors**

**propose a model of Dark Energy in which the field currently dominating (e) the energy density of the universe is (=) an "axion field" linearly coupled to (e&eb) the Pontryagin density, tr (FΛF), (i.e., the exterior derivative of (e) the Chern-Simons form) of (e) a massive gauge field. Authors assume that the axion has (e) self-interactions corresponding to (e&eb) a non-trivial (exponential) potential. Authors argue that a non-vanishing magnetic helicity of (e) the gauge field triggers (eb) slow-rolling of the axion at field values far below (eb) the Planck scale. Proposal leads to (eb) a "Tracking Dark Energy Scenario" in which the contribution of the axion energy density to (e) the total energy density is (=) constant (and small) during (e&eb) the early radiation phase, until (e) a secular growth term proportional to (e) the Pontryagin density of (e) the gauge field becomes (=) dominant. The initially small contribution of (e) the axion field to (e) the total energy density is related to the observed small baryon-to-entropy ratio. Some speculations concerning the nature of the gauge field are offered.**

**Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Cite as:arXiv: 1601.00057 [hep-th] (or arXiv:1601.00057v2 [hep-th] for this version) Tracking Dark Energy from Axion-Gauge Field Couplings Stephon Alexander, Robert Brandenberger, Juerg Froehlich (Submitted on 1 Jan 2016 (v1), last revised 23 Jan 2016 (this version, v2)) STOP Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=)**

**Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version) Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP In this work authors obtain bounds on (e&eb) the topological Abelian string-vortex and on (e&eb) the string-cigar, by using (e) a new measure of configurational complexity, known as (=) configurational entropy. In this way, the information-theoretical measure of (e) six-dimensional braneworld scenarios is capable to (e) probe situations where (e) the parameters responsible for (e) the brane thickness are arbitrary. The so-called configurational entropy (CE) selects (e&eb) the best value of the parameter in (eb) the model. This is accomplished by (e) minimizing the CE, namely, by (e) selecting the most appropriate parameters in the model that correspond to (e&eb) the most organized system, based upon (e) the Shannon information theory. This information-theoretical measure of complexity provides (eb) a complementary perspective to (e) situations where (e) strictly energy-based arguments are (=) inconclusive. Authors show that the higher the energy the higher the CE, what shows (eb) an important correlation between the energy of the localized field configuration and (e&eb) its associated entropic measure. Subjects: High Energy Physics - Theory (hep-th) DOI:10.1016/j.physletb.2016.02.038 Cite as: arXiv: 1601.00076 [hep-th] (or arXiv: 1601.00076v2 [hep-th] for this version) Bounds on topological Abelian string-vortex and string-cigar from information-entropic measure R. A. C. Correa, D. M. Dantas, C. A. S. Almeida, Roldao da Rocha (Submitted on 1 Jan**

**2016 (v1), last revised 17 Feb 2016 (this version, v2)) STOP Sum of the Holevo quantity (that bounds the capacity of quantum channels to transmit (e&eb) classical information about an observable) and (e&eb) the quantum discord (a measure of the quantumness of (e) correlations of that observable) yields (eb) an observable-independent total given by (e) the quantum mutual information This split naturally delineates ,dilates and disseminates information about (e&eb) quantum systems accessible to (e) observers -- information that is redundantly transmitted by (e&eb) the environment -- while showing that it is (=) maximized for (e) the quasi-classical pointer observable. Other observables are accessible only via (e&eb) correlations with the pointer observable .Further authors prove an anti-symmetry property relating accessible information and (e&eb) discord. It shows that information becomes (=) objective -- accessible to (e) many observers -- only as quantum information is relegated to (e) correlations with (e&eb) the global environment, (eb) and, therefore, locally inaccessible. The resulting complementarity explains (eb) why, in a quantum Universe, we perceive objective classical reality while (e) flagrantly quantum superpositions are out of reach. Subjects: Quantum Physics (quant-ph) Journal reference: Scientific Reports 3, 1729 (2013) Report number: LAUR-13-26386 Cite as:arXiv: 1303.4659 [quant-ph] (or arXiv: 1303.4659v2 [quant-ph] for this version) Complementarity of quantum discord and classically accessible information Michael Zwolak, Wojciech H. Zurek STOP Niels Bohr proposed that the outcome of the measurement becomes (=) objective and real, and, hence, classical, when (e) its results can be communicated by (e) classical means. In this work authors revisit Bohr's postulate using (e) modern tools from (e) the quantum information theory. They find a full confirmation of Bohr's idea: if a measurement device is in a nonclassical state, (eb) the measurement results cannot be communicated perfectly by classical means. In this case some part of information in the measurement apparatus is lost (e) in the process of communication: the amount of this lost information turns out to be (=) the quantum discord. The information loss occurs even when (e) the apparatus is not entangled**

**with the system of interest. The tools presented in this work allow (eb) to generalize Bohr's postulate: authors show that for pure system-apparatus states (eb) quantum communication does not provide any advantage when (e) measurement results are communicated to (e) more than one recipient. They further demonstrate the superiority of quantum communication to (e&eb) two recipients on (e&eb) a mixed system-apparatus state and show that (eb) this effect is fundamentally different from (e) quantum state cloning. Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. Lett. 111, 040401 (2013) DOI:10.1103/PhysRevLett.111.040401 Report number: LA-UR-12-25351 Cite as: arXiv: 1304.2914 [quant-ph] (or arXiv: 1304.2914v2 [quant-ph] for this version) Quantum discord cannot be shared Alexander Streltsov, Wojciech H. Zurek STOP Wojciech H. Zurek, Don N. Page investigate thermodynamic equilibrium of (e) a self-gravitating perfect fluid in (eb) a spherically symmetric system containing (e) a black hole of mass M by means of (e) the Tolman-Oppenheimer-Volkoff (TOV) equation. At r >> 2M its solutions describe (eb) a black-body radiation atmosphere with (e&eb) the Hawking temperature T_BH~1/(8 \pi M) that is increasingly blueshifted as r approaches (eb) 2M. However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments:**

**Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by (e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a firewall-like "Planck cocoon" with entropy close to (e) black hole entropy**

## Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

# Module Four

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equailibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-445) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Some of the concepts to which model is applicable is as follows:**

**Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by (e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave signal. The signal sweeps upwards in frequency from 35 to 250 Hz with (e&eb) a peak gravitational-wave strain of (e) 1.0×10−21. It matches the waveform predicted by general relativity for (e) the inspiral and merger of a pair of black holes and the ringdown of the resulting (eb) single black hole. The signal was observed with a matched-filter signal-to-noise ratio of (e) 24 and a false alarm rate estimated to be (=) less than 1 event per 203 000 years, equivalent to (=) a significance greater than 5.1σ. The source lies at (eb) a luminosity distance of 410+160−180 Mpc corresponding to (e&eb) a redshift z=0.09+0.03−0.04. In the source frame, the initial black hole masses are (=) 36+5−4M⊙ and 29+4−4M⊙, and the final black hole mass is (=) 62+4−4M⊙, with (e&eb) 3.0+0.5−0.5M⊙c2 radiated in (eb) gravitational waves. All uncertainties define (eb) 90% credible intervals. These observations demonstrate (eb) the existence of binary stellar-mass black hole systems. Observation of Gravitational Waves from a Binary Black Hole Merger B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 061102 – Published 11 February 2016 This is the first direct detection of gravitational waves and (e&eb) the first observation of a binary black hole**

**merger. Received 21 January 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.061102 STOP This is the second in a series of three papers on (e&eb) systematic analysis of rank 1 Coulomb branch geometries of (e) four dimensional N=2 SCFTs. In the first paper authors developed a strategy for (e) classifying physical rank-1 CB geometries of N=2 SCFTs. They show how to carry out this strategy computationally to construct (eb) the Seiberg-Witten curves and (e&eb) one-forms for all the rank-1 SCFTs. Explicit expressions are given for all cases, with the exception of the Nf=4 SU (2) gauge theory and (e&eb) the En SCFTs which were previously constructed. Classification includes all known rank-1 theories plus (+) a new one with (e&eb) an abelian flavor group, plus (+) nine additional theories whose existence is more speculative. Four of those, reported in first paper, depend on (=) the assumption of new frozen rank-1 SCFTs. Here authors also show that the assumption of the existence of certain rank-0 N=2 SCFTs leads to (eb) five additional consistent rank-1 CB geometries. Subjects: High Energy Physics - Theory (hep-th) Cite as:arXiv: 1601.00011 [hep-th] (or arXiv: 1601.00011v2 [hep-th] for this version Geometric constraints on the space of N=2 SCFTs II: Construction of special Kähler geometries and RG flows Philip C. Argyres, Matteo Lotito, Yongchao Lü, Mario Martone (Submitted on 31 Dec 2015 (v1), last revised 6 Jan 2016 (this version, v2)) STOP Authors propose a model of Dark Energy in which the field currently dominating (e) the energy density of the universe is (=) an "axion field" linearly coupled to (e&eb) the Pontryagin density, tr (F∧F), (i.e., the exterior derivative of (e) the Chern-Simons form) of (e) a massive gauge field. Authors assume that the axion has (e) self-interactions corresponding to (e&eb) a non-trivial (exponential) potential. Authors argue that a non-vanishing magnetic helicity of (e) the gauge field triggers (eb) slow-rolling of the axion at field values far below (eb) the Planck scale. Proposal leads to (eb) a "Tracking Dark Energy Scenario" in which the contribution of the axion energy density to (e) the total energy density is (=) constant (and small) during (e&eb) the early radiation phase, until (e) a secular growth term proportional to (e) the Pontryagin density of (e) the gauge field becomes (=) dominant. The initially small contribution of (e) the axion field to (e) the total energy density is related to the observed small baryon-to-entropy ratio. Some speculations concerning the nature of the gauge field are offered.**
**Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Cite as:arXiv: 1601.00057 [hep-th] (or arXiv:1601.00057v2 [hep-th] for this version) Tracking Dark Energy from Axion-Gauge Field Couplings Stephon Alexander, Robert Brandenberger, Juerg Froehlich (Submitted on 1 Jan 2016 (v1), last revised 23 Jan 2016 (this version, v2)) STOP Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version) Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP In this work authors obtain bounds on (e&eb) the topological Abelian string-vortex and on (e&eb) the string-cigar, by using (e) a new measure of configurational complexity, known as (=) configurational entropy. In this way, the information-theoretical measure of (e) six-dimensional braneworld scenarios is capable to (e) probe situations where (e) the parameters responsible for (e) the brane thickness are arbitrary. The so-called configurational entropy (CE) selects (e&eb) the best value of the parameter in (eb) the model. This is accomplished by (e) minimizing the CE, namely, by (e) selecting the most appropriate parameters in the model that correspond to (e&eb) the most organized system, based upon (e) the Shannon information theory. This information-theoretical measure of complexity provides (eb) a complementary perspective to (e) situations where (e) strictly energy-based arguments are (=) inconclusive. Authors show that the higher the energy the higher the CE, what shows (eb) an important correlation between the energy of the localized field configuration and (e&eb) its**

**associated entropic measure. Subjects: High Energy Physics - Theory (hep-th) DOI:10.1016/j.physletb.2016.02.038 Cite as: arXiv: 1601.00076 [hep-th] (or arXiv: 1601.00076v2 [hep-th] for this version) Bounds on topological Abelian string-vortex and string-cigar from information-entropic measure R. A. C. Correa, D. M. Dantas, C. A. S. Almeida, Roldao da Rocha (Submitted on 1 Jan 2016 (v1), last revised 17 Feb 2016 (this version, v2)) STOP Sum of the Holevo quantity (that bounds the capacity of quantum channels to transmit (e&eb) classical information about an observable) and (e&eb) the quantum discord (a measure of the quantumness of (e) correlations of that observable) yields (eb) an observable-independent total given by (e) the quantum mutual information This split naturally delineates ,dilates and disseminates information about (e&eb) quantum systems accessible to (e) observers -- information that is redundantly transmitted by (e&eb) the environment -- while showing that it is (=) maximized for (e) the quasi-classical pointer observable. Other observables are accessible only via (e&eb) correlations with the pointer observable .Further authors prove an anti-symmetry property relating accessible information and (e&eb) discord. It shows that information becomes (=) objective -- accessible to (e) many observers -- only as quantum information is relegated to (e) correlations with (e&eb) the global environment, (eb) and, therefore, locally inaccessible. The resulting complementarity explains (eb) why, in a quantum Universe, we perceive objective classical reality while (e) flagrantly quantum superpositions are out of reach. Subjects: Quantum Physics (quant-ph) Journal reference: Scientific Reports 3, 1729 (2013) Report number: LAUR-13-26386 Cite as:arXiv: 1303.4659 [quant-ph] (or arXiv: 1303.4659v2 [quant-ph] for this version) Complementarity of quantum discord and classically accessible information Michael Zwolak, Wojciech H. Zurek STOP Niels Bohr proposed that the outcome of the measurement becomes (=) objective and real, and, hence, classical, when (e) its results can be communicated by (e) classical means. In this work authors revisit Bohr's postulate using (e) modern tools from (e) the quantum information theory. They find a full confirmation of Bohr's idea: if a measurement device is in a nonclassical state, (eb) the measurement results cannot be communicated perfectly by classical means. In this case some part of information in the measurement apparatus is lost (e) in the process of communication: the amount of this lost information turns out to be (=) the quantum discord. The information loss occurs even when (e) the apparatus is not entangled with the system of interest. The tools presented in this work allow (eb) to generalize Bohr's postulate: authors show that for pure system-apparatus states (eb) quantum communication does not provide any advantage when (e) measurement results are communicated to (e) more than one recipient. They further demonstrate the superiority of quantum communication to (e&eb) two recipients on (e&eb) a mixed system-apparatus state and show that (eb) this effect is fundamentally different from (e) quantum state cloning. Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. Lett. 111, 040401 (2013) DOI:10.1103/PhysRevLett.111.040401 Report number: LA-UR-12-25351 Cite as: arXiv: 1304.2914 [quant-ph] (or arXiv: 1304.2914v2 [quant-ph] for this version) Quantum discord cannot be shared Alexander Streltsov, Wojciech H. Zurek STOP Wojciech H. Zurek, Don N. Page investigate thermodynamic equilibrium of (e) a self-gravitating perfect fluid in (eb) a spherically symmetric system containing (e) a black hole of mass M by means of (e) the Tolman-Oppenheimer-Volkoff (TOV) equation. At r >> 2M its solutions describe (eb) a black-body radiation atmosphere with (e&eb) the Hawking temperature T_BH~1/(8 \pi M) that is increasingly blueshifted as r approaches (eb) 2M. However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments: Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by (e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a firewall-like "Planck cocoon" with entropy close to (e) black hole entropy. Authors reproduce paper with an updated title and abstract. For a later study with similar results see G. 't Hooft, gr-qc/9706058 Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Physical Review D, Volume 29, Number 4, pp. 628 - 631 (1984) DOI:10.1103/PhysRevD.29.628 Cite as:arXiv:1511.07051 [gr-qc] (or arXiv:1511.07051v1 [gr-qc] for this version) Firewalls, black-hole thermodynamics, and singular solutions of the Tolman-Oppenheimer-Volkoff equation Wojciech H. Zurek, Don N. Page (Submitted on 22 Nov 2015) STOP Frederic Magniez, Ashwin Nayak consider the problem of testing the commutativity of (e) a black-box group specified by its k generators. The complexity (in terms of k) of (e) this problem was first considered by Pak, who gave (eb) a**

**randomized algorithm involving O(k) group operations. They construct a quite optimal quantum algorithm for this problem whose complexity is in O~(k2/3)O~(k2/3). The algorithm uses and highlights the power of the quantization method of Szegedy. For the lower bound of Ω(k2/3)Ω(k2/3), authors give a reduction from (e) a special case of Element Distinctness to our problem. Along the way, they prove the optimality of the algorithm of (e) Pak for the randomized model. Algorithmica July 2007, Volume 48, Issue 3, pp 221-232 First online: 25 June 2007 Quantum Complexity of Testing Group Commutativity Frederic Magniez, Ashwin Nayak STOP Pawel Wocjan, Jon Yard analyze relationships between quantum computation and (e&eb) a family of generalizations of the Jones polynomial. Extending recent work by Aharonov et al., authors give efficient quantum circuits for (e) implementing the unitary Jones-Wenzl representations of the braid group. We use these to provide new quantum algorithms for (e) approximately evaluating a family of specializations of the HOMFLYPT two-variable polynomial of trace closures of braids. They also give (eb) algorithms for approximating the Jones polynomial of a general class of closures of (e) braids at roots of unity. Next authors provide (eb) a self-contained proof of a result of (e) Freedman et al. that any quantum computation can be replaced by (e&eb) an additive approximation of the Jones polynomial, evaluated at (eb) almost any primitive root of unity. Proof encodes two-qubit unitaries into (e&eb) the rectangular representation of (e) the eight-strand braid group. Authors then give QCMA-complete and (e&eb) PSPACE-complete problems which are based on braids. They conclude with direct proofs that evaluating the Jones polynomial of the plat closure at most primitive roots of unity is a #P-hard problem, while learning its most significant bit is PP-hard, circumventing (e) the usual route through the Tutte polynomial and graph coloring. Subjects: Quantum Physics (quant-ph); Geometric Topology (math.GT) Cite as: arXiv:quant-ph/0603069 (or arXiv:quant-ph/0603069v3 for this version) The Jones polynomial: quantum algorithms and applications in quantum complexity theory Pawel Wocjan, Jon Yard STOP Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan address the task of switching off the Hamiltonian of (e) a system by removing (e) all internal and system-environment couplings. Authors propose dynamical decoupling schemes, which use (e) only bounded-strength controls, for quantum many-body systems with local system Hamiltonians and local environmental couplings. To do so, they introduce (eb) the combinatorial concept of (e) balanced-cycle orthogonal arrays (BOAs) and show how to construct them from (e) classical error-correcting codes. The derived decoupling schemes may be useful as (=) a primitive for more complex schemes, e.g., for Hamiltonian simulation. For the case of n qubits and a 2-local Hamiltonian, the length of the resulting decoupling (e&eb) scheme scales as O (nlogn), improving over (e&eb) the previously best-known schemes that scaled quadratically with (e&eb) n. More generally, using balanced-cycle orthogonal arrays constructed from families of BCH codes, we show that bounded-strength decoupling for any ℓ-local Hamiltonian, where ℓ≥2, can be achieved using decoupling schemes of length at most O(nℓ−1logn). Subjects:Quantum Physics (quant-ph); Emerging Technologies (cs.ET) Journal reference: IEEE Trans. Inf. Theory, 62(5): 2881-2894, 2016 DOI: 10.1109/TIT.2016.2535183 Report number: MIT-CTP /4705 Cite as: arXiv: 1509.00408 [quant-ph] (or arXiv: 1509.00408v2 [quant-ph] for this version) Improved bounded-strength decoupling schemes for local Hamiltonians Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan STOP Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola propose dynamical control schemes for Hamiltonian simulation in many-body quantum systems that avoid instantaneous control operations and rely solely on (e) realistic bounded-strength control Hamiltonians. Each simulation protocol consists of (e) periodic repetitions of a basic control block, constructed as (=) a suitable modification of an "Eulerian decoupling cycle," that would otherwise implement (eb) a trivial (zero) target Hamiltonian. For an open quantum system coupled to (e) an uncontrollable environment, approach may be employed to (e) engineer an effective evolution that simulates (e&eb) a target Hamiltonian on the system, while suppressing (e) unwanted decoherence to the leading order. Authors present illustrative applications to (e) both closed- and open-system simulation settings, with (e&eb) emphasis on simulation of non-local (two-body) Hamiltonians using (e) only local (one-body) controls. In particular, authors provide (eb) simulation schemes applicable to (e&eb) Heisenberg-coupled spin chains exposed to (e) general linear decoherence, and show how to simulate (e&eb) Kitaev's honeycomb lattice Hamiltonian starting from Ising-coupled qubits, as (=) potentially relevant to (e&eb) the dynamical generation of a topologically protected (e&eb) quantum memory. Additional implications for quantum information processing are discussed. Subjects: Quantum Physics (quant-ph) Journal reference: New J. Phys. 16, 045021 (2014) DOI: 10.1088/1367-2630/16/4/045021 Cite as:arXiv: 1310.4153 [quant-ph] (or arXiv: 1310.4153v1 [quant-ph] for this version) Hamiltonian quantum simulation with bounded-strength controls Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola STOP Felix Fontein, Pawel Wocjan**

**present a quantum algorithm for computing the period lattice of (e) infrastructures of fixed dimension. The algorithm applies to (e&eb) infrastructures that satisfy certain conditions. The latter are always fulfilled for (e) infrastructures obtained from (e) global fields, i.e., algebraic number fields and function fields with (e&eb) finite constant fields. The first main contributions is an exponentially better method for (e) sampling approximations of vectors of the dual lattice of (e) the period lattice than the methods outlined in the works of (e) Hallgren and Schmidt and Vollmer. This new method improves (eb) the success probability by (e) a factor of at least 2^ {n^2-1} where (e) n is the dimension. The second main contribution is a rigorous and complete proof that the running time of the algorithm is (=) polynomial in the logarithm of (e) the determinant of the period lattice and (e&eb) exponential in n. The third contribution is the determination of (e) an explicit lower bound on (e&eb) the success probability of our algorithm which greatly improves on (e&eb) the bounds given in the above works. The exponential scaling seems inevitable because (e) the best currently known methods for (e) carrying out fundamental arithmetic operations in infrastructures obtained from (e) algebraic number fields take exponential time. In contrast, the problem of computing (e&eb) the period lattice of infrastructures arising from function fields can be solved without (e) the exponential dependence on (e&eb) the dimension n since this problem reduces efficiently to the abelian hidden subgroup problem. This is also true for (e) other important computational problems in (eb) algebraic geometry. The running time of the best classical algorithms for (e) infrastructures arising from (e) global fields increases subexponentially with (e&eb) the determinant of the period lattice. Subjects: Quantum Physics (quant-ph) Cite as: arXiv: 1111.1348 [quant-ph] (or arXiv: 1111.1348v2 [quant-ph] for this version) Quantum Algorithm for Computing the Period Lattice of an Infrastructure Felix Fontein, Pawel Wocjan STOP Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan advocate a new approach of (e) addressing hidden structure problems and (e&eb) finding efficient quantum algorithms. Authors introduce and investigate the Hidden Symmetry Subgroup Problem (HSSP), which is a generalization of the well-studied Hidden Subgroup Problem (HSP). Given a group acting on (e&eb) a set and an oracle whose level sets define (eb) a partition of the set; the task is to recover (e) the subgroup of symmetries of (e) this partition inside the group. The HSSP provides (eb) a unifying framework that, besides the HSP, encompasses (e) a wide range of algebraic oracle problems, including (e) quadratic hidden polynomial problems. While the HSSP can have (e) provably exponential quantum query complexity, authors obtain efficient quantum algorithms for various interesting cases. To achieve this, authors present a general method for (e) reducing the HSSP to the HSP, which works efficiently in several cases related to (e&eb) symmetries of polynomials. The HSSP therefore connects in a rather surprising way certain hidden polynomial problems with (e&eb) the HSP. Using this connection, authors obtain the first efficient quantum algorithm for (e) the hidden polynomial problem for (e) multivariate quadratic polynomials over fields of (e) constant characteristic. Authors also apply the new methods to (e&eb) polynomial function graph problems and (e&eb) present an efficient quantum procedure for (e) constant degree multivariate polynomials over (e&eb) any field. This result improves in several ways the currently known algorithms. Subjects: Quantum Physics (quant-ph) Journal reference: SIAM Journal on Computing 42:5 (2013) DOI: 10.1137/120864416 Cite as: arXiv: 1107.2189 [quant-ph] (or arXiv: 1107.2189v2 [quant-ph] for this version) Hidden Symmetry Subgroup Problems Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan STOP Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan present an efficient general method for realizing a quantum walk operator corresponding to (e&eb) an arbitrary sparse classical random walk. Approach is based on Grover and Rudolph's method for preparing coherent versions of (e) efficiently integrable probability distributions. This method is intended for (e) use in quantum walk algorithms with polynomial speedups, whose complexity is usually measured in terms of (e&eb) how many times we have to apply a step of a quantum walk, compared to (e&eb) the number of necessary classical Markov chain steps. Authors consider a finer notion of complexity including (e) the number of elementary gates it takes to implement each step of the quantum walk with (e&eb) some desired accuracy. The difference in complexity for (e) various implementation approaches is that method scales linearly in the sparsity parameter and (e&eb) poly-logarithmically with (e&eb) the inverse of the desired precision. Best previously known general methods either scale quadratically in (eb) the sparsity parameter, or polynomially in (eb) the inverse precision. Approach is especially relevant for (e) implementing quantum walks corresponding to (e&eb) classical random walks like those used in (eb) the classical algorithms for approximating permanents and sampling from binary contingency tables. In those algorithms, the sparsity parameter grows with (e&eb) the problem size, while (e) maintaining high precision is required. Subjects: Quantum Physics (quant-ph) Journal reference: Quantum Information and Computation, Vol.10,**

**No.5&6, pp0420-0434 (2010) Cite as: arXiv: 0903.3465 [quant-ph] (or arXiv: 0903.3465v3 [quant-ph] for this version) Efficient Circuits for Quantum Walks Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan STOP Preparing the ground state of a system of interacting classical particles is (=) an NP-hard problem. Thus, there is in general no better algorithm to (e) solve this problem than exhaustively going through all N configurations of (e) the system to determine (eb) the one with lowest energy, requiring (e) a running time proportional to N. A quantum computer, if it could be built, could solve (e) this problem in time sqrt (N). Here, authors present a powerful extension of this result to (e) the case of interacting quantum particles, demonstrating (eb)that a quantum computer can prepare the ground state of (e) a quantum system as efficiently as it does for (e) classical systems. Subjects: Quantum Physics (quant-ph) Journal reference: Phys.Rev.Lett.102:130503, 2009 DOI:10.1103/PhysRevLett.102.130503 Cite as: arXiv: 0809.2705 [quant-ph] (or arXiv: 0809.2705v1 [quant-ph] for this version) Preparing ground states of quantum many-body systems on a quantum computer David Poulin, Pawel Wocjan STOP Daniel Nagaj, Pawel Wocjan construct a simple translationally invariant, nearest-neighbor Hamiltonian on (e&eb) a chain of 10-dimensional qudits that makes it possible to (e) realize universal quantum computing without (e) any external control during (e&eb) the computational process. Authors only require the ability to prepare an initial computational basis state which encodes (e) both the quantum circuit and its input. The computational process is (=) then carried out by (e) the autonomous Hamiltonian time evolution. After a time polynomially long in the size of (e) the quantum circuit has passed, the result of (e) the computation is obtained with (e&eb) high probability by (e) measuring a few qudits in the computational basis. This result also implies (eb) that there cannot (e) exist efficient classical simulation methods for (e) generic translationally invariant nearest-neighbor Hamiltonians on (e&eb) qudit chains, unless (e) quantum computers can be efficiently simulated by (e&eb) classical computers (or, put in complexity theoretic terms, unless BPP=BQP). Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. A 78, 032311 (2008) DOI: 10.1103/PhysRevA.78.032311 Cite as: arXiv: 0802.0886 [quant-ph] (or arXiv: 0802.0886v4 [quant-ph] for this version) Hamiltonian Quantum Cellular Automata in 1D Daniel Nagaj, Pawel Wocjan**

## 10. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_4}{dt} = -b_4 Q_4 + g_4(g_4)Q_4$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_4(0) = 0, \quad g_4(P_4) > 0, \ P_4 \ 0, \qquad \lim_{P_4 \to \infty} g_4(P_4) = (K_4)_4 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_4 + (\rho_a)_4 + D_4(a_4, P_4, Q_4)\rho_4 + 0, \qquad 0 < a_4 ?\infty,$$

$$\rho_4(t,0) = \int_0^{\infty} n_4(a_4, P_4, Q_4)\rho_4(t_4, a_4)da_4$$

Where $D_4$ is the mortality function and $n_4$ is the fertility function.

Put $_4$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_4}{dt}\right]_4 = x_4\left[\frac{dQ_2}{dt}\right]_4 - x_2\left[\frac{dQ_4}{dt}\right]_4$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_4 - x_4\left[\frac{dQ_2}{dt}\right]_4$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_4 - x_6\left[\frac{dQ_3}{dt}\right]_4$$

Where $x_4, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_4$ and $g_4(P_4)Q_4$

Now when you add the term $(b_4Q_4)$ with $g_4(P_4)(Q_4)$ we are infect see is that $(b_4Q_4)$ is credited incrementally part of $g_4(P_4)(Q_4)$ or instant nearly, when subtracted from $\frac{dQ_4}{dt}$ , it means that $g_4(P_4)(Q_4)$ is slowly eaten by $\frac{dQ_4}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_4 = \mu_4(a_4) + (f_4)_4(P_4)_4 + (f_2)_4(Q)_4 \quad (2)$$

$$n_4 = \beta_4(a_4)h_4(Q_4), \quad (3)$$

Where $\mu_4$the age- dependent debit rate whit is $\mu_4(a_4)$ ⊠0 $\forall a > 0$; $(f_4)_4$ is the density- dependent death rate with $f_4(0) = 0$, $(f_4^4)_4(P_4)0$, $\forall P$ ⊠0, and $\lim_{p\to\infty}(f_4)_4(P_4) = (k_4)_4$ ⊠∞; $f_2$ is the functional response of creditors to prey with $[f_2]_4(0) = 0, (f_2^4)_4(Q) > 0, \forall Q$ ⊠0, and $\lim_{Q\to\infty}(f_2)_4(Q_4) = [k_2]$ ⊠∞; $\beta$ is the age dependent birth rate with $\beta(a)$ ⊠0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 4, h^4(Q)$ ⊠0, $\forall Q$ ⊠0, and $\lim_{Q\to\infty} h(Q) = k_3$ ⊠0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_4 + (\rho_a)_4 = -\mu_4(a_4) + (f_4)_4P_4) + (f_2)_4(Q_4))(\rho_4)_4 \quad (4)$$

$$\rho_4(t_4, 0) = h_4(Q_4)\int_0^\infty \beta_4(a_4)\rho_4(t, a_4)da, \quad (5)$$

$$\frac{dQ_4}{dt} = -b_4Q_4 + g_4(P_4)Q_4 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_4$ and $Q_4$ of (4-6) satisfy

$$0 \leq \rho_4(t, a_4) \leq \rho_4^*(t, a_4),$$

$$0 \leq Q_4(t) \leq Q_4^*(t),$$

Where $\rho_4^*(t, a_4)$ and $Q_4^*(t)$, are solutions of

$$(\rho_t^*)_4 + (\rho_a^*)_4 = -(\mu_4(a_4) + (f_4)_4(P_4^*) + (f_2)_4Q_4^*))\,\rho_4^*, \quad (7)$$

$$\rho_4^*(t,0) = \int_0^\infty \beta_4(a_4)\,\rho_4^*(t,a_4)da_4 \quad (8)$$

$$\frac{dQ_4^*}{dt} = -b_4\,Q_4^* + g_4(P_4^*)Q_4^*, \quad (9)$$

With $P_4^*(t) = \int_0^\infty \rho_4^*(t,a_4)da_4$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_4^*(t)}{dt} = P_4^*(t)\big(r_4 - (f_4)(P_4^*) - (f_2)_4(Q^*)\big), \quad (10)$$

$$\frac{dQ_4^*(t)}{dt} = Q_4^*(t)(-b_4 + g_4(P_4^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_4(a)e^{-\int_0^a \mu_4(\tau)dt}e^{-ra}\,da_4 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_4^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(t)dt}\,da_4 > 1, \quad (H1)$$

Which guarantees $r_4 > 0$.

*Remark.* If $\beta_4(a) = [\beta_0]_4 a_4 e^{-\alpha_4 a_4}, \mu_4(a_4) \equiv \mu_4$ then the solution of (12) is

$$r_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_4]_4(P_4) = 0, [f_2](Q_4) = \lambda_4 Q_4\ h_4(Q_4) = 1, \qquad and\ \ g_4(P_4) = c_4 P_4$$

$$[\rho_4]_4 + [\rho_a]_4 = -(\mu_4 + \lambda_4 Q_4)\rho_4$$

$$\rho_4(t,0) = \int_0^\infty [\beta_0]_4 a_4 e^{-\alpha_4 a_4} \rho_4(t,a_4)\, da_4$$

$$\frac{dQ_4}{dt} = -b_4 Q_4 + c_4 P_4 Q_4$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_4}{dt} = \omega_4 P_4 - \lambda_4 P_4 Q_4$$

$$\frac{dQ_4}{dt} = b_4 Q_4 - c_4 P_4 Q_4$$

Where $\omega_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$

## 11. THE EQUILIBRIUM $(\rho_0)_4(a_4), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_4(a_4), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_4 + (\rho_a)_4 = -\big(\mu_4(a_4) + (f_4)_4(P_4)\big)\rho,$$

$$\rho_4(t,0) = \int_0^\infty \beta_4(a_4)\rho_4(t,a_4)da_4 \qquad (13)$$

The solution of (13) has the form of

$$\rho_4(t,a_4) = Ce_4^{r(t-a_4)} - \int_0^a (\tau)d\tau - \int_0^4 (f_4)_4\big(P_4(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_4(t) = C_4 e^{r_4 t - \int_0^t (f_4)_4(P_4(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_4^{-4}(r)$ if $k_4 > r$. This is equivalent to

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{(k_4)_4 a_4} da_4 < 1.$$

In fact, if we define

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau)d\tau} e^{-s_4 a_4}\, da_4$$

$R(s)$ is a decreasing function. $R_4(r_4) = 1$ And $R_4(k_4)_4 < 1$ imply $r_4 < (k_4)_4$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_4(a_4), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{k_4 a_4} da_4$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_4 \equiv \int_0^\infty (\rho_0)_{44}(a_4)da_4$ *satisfies*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{-(f_4)_4(P_0)a_4} da$$

$$= 1 \qquad (11)$$

Next, we study the stability of $(\rho_0)_4(a_4), 0)$

Let $u_4(t, a_4) = \rho(t, a_4) - \rho_0(a_4)$ Then

$$N_4(t) \equiv \int_0^\infty u_4(t, a_4)da_4 = \int_0^\infty \rho_4(t, a_4)da_4 - \int_0^\infty (\rho_0)_4(a_4)da = P_4(t) - P_0)_4$$

And

$$(u_4)_4 + (u_a)_4 = (\rho_t)_4 + (\rho_a)_4 - (\rho_0^4)_4(a_4)$$

$$= -\mu_4(a_4) + (f_4)_4(P_4) + (f_2)_4(Q_4))(u_4 + (\rho_0)_4) + \mu_4(a_4)(\rho_0)_4 + (f_4)_4(P_0)_4(\rho_0)_4$$

$$= -\mu_4(a_4)u_4 - (f_4)_4(P_4) + (f_2)_4(Q_4))u_4 - (\rho_0)_4(f_4)_4(P_4) - (f_4)_4(P_0)_4$$

$$+ (f_2)_4(Q_4)), \qquad (15)$$

Where $(\rho_0^4)_4(a_4)$ denotes $d(\rho_0)_4(a_4)/da_4$

Expansion $(f_4)_4(P_4)$ and $(f_2)_4(Q_4)$ at $(P_0)_4$ and 0 yield respectively

$$(f_4)_4(P_4) = (f_4)_4((P_0)_4 + N_4)$$

$$= (f_4)_4(P_0)_4 + (f_4^4)_4(P_0)_4 N_4 + o_4(N_4), \qquad (16)$$

$$(f_2)_4(Q_4)$$

$$= (f_2^4)_4(0)Q + o_4(Q_4) \quad (17)$$

Write $f_4(P_0) = f_{40}, f_4^{'}(P_0) = f_{40}^{'}$, and $f_2^{'}(0) = f_{20}^{'}$ and substitute (16) and (17) into (15) to obtain

$$(u_4)_4 + (u_a)_4 \approx -(\mu_4(a_4) + (f_{40})_4)u - (f_{40}^4)_4(\rho_0)_4(N_4) - (f_{20}^4)_4(\rho_0)_4(Q_4)$$

$$u_4(t,0) = \rho_4(t,0)(\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a) + (\rho_0)_4(a_4)da_4 - (\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a_4)da_4 + (h_4(Q_4) - 4)(\rho_0)_4(0) \quad (18)$$

Expanding

$$h_4(Q_4) = 1 + (h^4)_4(0)Q_4 + 0(Q_4) \equiv 1 + h_0^{'}Q_4 + 0Q_4$$

And substituting it into (48), we arrive at

$$u_4(t,0) \approx \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + (h_{a_4}^{'})(\rho_0)_4Q_4$$

The linearization is the following :

$$[u_4] + [u_4] = -(\mu_4(a_4) + [t_{40}]_4u_4 - [\alpha_4]_4(a_4)N_4 - [\alpha_2]_4(a_4)Q_4$$

$$u_4(t,0) = \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + [\xi_4]Q_4 \quad (19)$$

$$\frac{dQ_4}{dt} = -b_4Q_4$$

where

$$[\alpha_4]_4(a_4) = [t_4^4 a]_4[p, o]_4(a_4)d2(a_4) = [f_{20}^4]_4(\rho\, o\,)_4(a_4)\xi = [\mathrm{h\ o}^4]_4[\rho \mathrm{o}]_4(\mathrm{o}) \text{ and } \mathrm{b}_4 = \mathrm{b}_4 - g_4(P_4)$$

If $g_4(P_0) > b$, clearly, $(\rho_0)_4(a_4), 0)$ is unstable and the predator population is persistent. Assume $g_4(P_0)_4 < b$.

Then, $Q_4(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_4 + u_a = -(\mu_4(a_4) + (f_{40})_4)u_4 - (\alpha_4)_4(a_4)N_4$$

$$= -(\mu_4(a_4) + f_{40})u_4 - \alpha_4(a_4)\int_0^\infty u_4(t, a_4)da_4$$

$$u_4(t, 0) = \int_0^\infty \beta_4\,(a_4)u_4(t, a_4)da_4$$

Using the method developed by Cushing [2, 3], we suppose

$$u_4(t, a_4) = \tilde{\mathrm{u}}_4(a_4)^{e^{\gamma(t-a_4)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_4}\, da_4$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_4)$ gives,

$$\tilde{\mathrm{u}}_4(a_4) = e^{-\int_0^a \mu_4(\tau_4)d\tau - (f_{40})_4 a}\left(\tilde{\mathrm{u}}_4(0) - w\int_0^a e^{\int_0^\tau \mu_4(s_4)ds + (f_{40})}\tau_4\,(\alpha_4)_4(\tau_4)e^{\gamma_4\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_4(a_4)e^{-\gamma_4 a_4 - (f_{40})_4 a_4 - \int_0^a \mu_4(\tau_4)d\tau}\right)\tilde{\mathrm{u}}_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a(\alpha_4)_4(\tau_4)e^{-\gamma_4\tau_4}e^{(f_{40})_4\tau_4+\int_0^{\tau}\mu_4(s_4)ds}d\tau_4\,da_4$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}da\right)(\tilde{u}_4(0))_4$$

$$+\left(-4-\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a e^{(f_{40})_4\tau_4+\int_0^{\tau}\mu_4(s_4)ds+y\tau}\alpha_4(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(t_4)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_4\,(a_4)Y_4(\gamma_4,a_4)da_4\right)u_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)\int_0^a(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4=0,$$

(24)

$$\left(\int_0^{\infty}Y_4(\gamma_4,a_4)da_4\right)\tilde{u}_4(0)$$

$$+\left(-1-\int_0^{\infty}Y_4(\gamma_4,a_4)\int_0^a(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4\right)w=0.$$

Equqtion (20) has a solution of the form of $u_4(t,a_4) = u_4(a_4)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_4(0) \not\equiv 0$ and $\gamma, w$ in the complwx plane.

Let $I_4(\gamma_4,a_4) = \int_0^a (\rho_0)_4(\tau_4)Y^{-4}(\gamma_4,a_4)d\tau_4$. We obtain the characteristic equation

$$f_{40}^{l}{}_4\left(\int_0^\infty Y_4(\gamma_4,a_4)da_4\right)\left(\int_0^\infty \beta_4(a_4)Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$+\left(1-\int_0^\infty \beta_4(a_4)Y_4(\gamma_4,a_4)da\right)\left(1+(f_{40}^{l})_4\int_0^\infty Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$= 0. \qquad (25)$$

Since

$$(\rho_0)_4(a_4) = (\rho_0)_4(0)e^{-\int_0^a \mu_4(\tau_4)d\tau_4-(f_{40})_4a_4}$$

$$I_4(\gamma_4,a_4) = \int_0^a (\rho_0)_4(\tau_4)e^{\gamma_4\tau_4+(f_{40})_4\tau+\int_0^\tau \mu_4(s_4)ds_4}d\tau$$

$$= \frac{(\rho_0)_4(0)}{\gamma_4}(e^{\gamma_4a_4}-1)$$

Let

$$E_4(a_4) = e^{-(f_{40})_4a_4-\int_0^a \mu_4(\tau_4)d\tau_4}$$

Then $[Y_4]_4[\gamma_4]_4,a_4) = E_4(a_4)e^{-\gamma_4a_4}$. by a simple calculation,

$$\int_0^\infty \beta_4(a_4)Y_4(\gamma_4,a_4)I(\gamma_4,a_4)da = \frac{[\rho_0]_4(0)}{\gamma}\left(1-\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}\ da\right);$$

$$\int_0^\infty Y_4(\gamma_4,a_4)I(\gamma_4,a_4)da = \frac{[\rho_0]_4(0)}{\gamma_4}\left(\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}\ da_4\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_4(0)}{\gamma_4}[f_{40}^4]\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}da_4\ (\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4)$$

$$+\left(1-\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right).$$

$$\left(1+\frac{(\rho_0)_4(0)}{\gamma_4}(f_{40}^{'})_4\left(\int_0^\infty E_4(\,a_4)da_4-\int_0^\infty E_4(\,a_4)e^{-\gamma_4 a_4}da_4\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right)(1+\frac{(f_{40}^{'})_4}{\gamma_4}(\rho_0)_4(0)\int_0^\infty E_4(\,a_4)da_4)=0.$$

We need ũ(a) $\not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)da_4 = 1-\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4 \neq 0,$$

Since $\int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)\ \int_0^a \alpha_4\,(\tau_4)Y^{-4}(\gamma_4,\tau_4)d\tau_4 da_4 > 0$. Therefore, the characteristic solution is

$$\gamma_4 = -(f_{40}^{'})_4(\rho_0)_4 \int_0^{\infty} E_4(a_4)da_4.$$

Since $f_4'(P_0)_4 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_4), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_4), 0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_4 + (\mu_4(a_4) + [f_4]_4(P_4^*) + [f_2]_4(Q_4^*))P_4^* = 0, \quad (26)$$

$$p_4^*(0) = h_4(Q_4^*)\int_0^{\infty} \beta_4(a_4)p_4^*(a_J)p_4^*(a_4)da_4 \quad (27)$$

$$g_4(P_4^*) - b_4 = 0, \quad (28)$$

Where $P_4^* = \int_0^{\infty} p_4^*(a_4)\, da_4$

Equation (28) yields $P^* = g^{-4}(b)$,

From (26),

$p_4^*(a) = p_4^*(0)e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}$

Substituting into (27), we need to solve

$$\left(1 - h_4(Q_4^*)\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}da_4\right)p_4^*(0) = 0$$

For $Q_4^*$

Fix $p_4^*(0)$ and $P_4^*$ Let

$$H_4(z_4) = 1 - h_4(z_4)\int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}da_{,4}$$

And assume that

$$[k_3]_4\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(g_4^{-4}(b_4))a_4 - [k_2]_4 a_4}\, da_4 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_4(z_4)] > 0.$$

Clearly,

$$\frac{dH_4}{dz_4} = h_4(z)([f_2']_4(z_4))\int_0^\infty a_4\beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - f_4(P_4^*)a_4 - f_2(z_4)a_4}da_{\square}$$

$-h'(z_4)\int_0^\infty \beta_4(a_4)\, e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}\quad da_4$

the assumptions on $\left[f_2^4\right]_4 and\ [h^4]^4 dH_4/dz > 0$. Hence, if

$$H_4(0) = 1 - \int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4} da_4 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_4(0) < 0$ follows if $g_4^-(b_4) < [f_4^{-4}]_4(r_4)$, where $r_4$ is the solution of (12).

Indeed, since

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 e^{-y\frac{a}{4}} da_4}$$

Is a decreasing function of $s_4$ when $P_4^* = [g^{-4}]_4(b_4) < [f_4^{-4}]_4(r_4)$, that is $[f_4]_4(P_4^*) < r$, $R\left(f_4(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_4(g^{-4}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4\left(g_4^{-4}(b)\right)a_4} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_4^*(a_4), Q_4^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_4^*(a_4), Q_4^*)$ is more complicated. We again consider perturbations $m_4(t_4 a_4) = p_4(t_4, a_4) - p_4^*(a_4)$ and $W_4(t) = Q_4(t) - Q^*$.Then

$$M_4(t) \equiv \int_0^\infty m_4(t, a_4) da_4 = P_4(t) - P_4^*,$$

$$[m_4]_4 + [m_a]_4 = [p_4]_4 + [p_a]_4 - [p^*]_4(a_4) \approx -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4$$

$$-[f_{4*}^4]_4 p_4^*(a_4)M - [f_{2*}^4]_4 p_4^*(\mathrm{a})W_4,$$

$$m_4(t,0)p_4(t,0) - p_4^*(0) \approx h_* \int_0^\infty \beta_4(a_4)m_4(t,a_4)\,da + \frac{[h_*^4]_4}{[h_*]_4}\, p_4^*(0)W_4$$

$$\frac{dW_4}{dt} = \frac{dQ_4}{dt} \approx [g_*^4]_4 Q_4^* M_{,4}$$

Where $[f_{4*}]_4$ denoting $[f_4(P^*)]_4$, $[f_{4*}]_4$ denoting $[f_2(P^*)]_4$, $[f_{4*}^{'}]_4$ denoting $[f^{'}{}_4(P^*)]_4 [f_{2*}^{'}]_4$ denoting $[f_2]_4$ $(Q_4^*,$ $h_{*}$ denoting $h_*^{'}$ denoting $h_4^{'}$ $(Q_4^*)$, and $[g_*^{'}]_4$

Denoting $[g^{'}]_4(P_4^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_4]_4 + [m_a]_4 = -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4 - [f^{'}{}_{4*}]_4 [p_4^*](a_4)M_4 - [f_{2*}^{'}]_4 p_4^*(a_4)W_4$$

$$m_4(t,0) = [h_*]_4 \int_0^\infty \beta_4(a_4)m_4(t,a_4)da_4 + \frac{[h_*^4]_4}{[h^*]_4}[p^*]_4(0)\, W_{,4} \qquad (30)$$

$$\frac{dW_4}{dt} = [g_*^4]_4 [Q_4^*]M_4$$

This is equation is important in the following cases: problem of using observed correlations to infer causal relations is relevant to (e) a wide variety of scientific disciplines. Yet given correlations between just two classical variables, (eb) it is impossible to determine whether they arose from (e) a causal influence of one on the other or a common cause influencing both, unless one can implement (eb) a randomized intervention. **Katja Ried, Megan Agnew,**

**Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. Resch** here consider the problem of causal inference for (e) quantum variables. Authors introduce causal tomography, which unifies and generalizes (e&eb) conventional quantum tomography schemes to provide (eb) a complete solution to the causal inference problem using (e) a quantum analogue of a randomized trial. Authors furthermore show that, in contrast to the classical case, observed quantum correlations alone can sometimes provide (eb) a solution. They implement a quantum-optical experiment that allows (eb) to control (e&eb) the causal relation between two optical modes, and two measurement schemes -- one with (e&eb) and one without (e) randomization -- that extract this relation from (e) the observed correlations. Results show that entanglement and (e&eb) coherence, known to be central to (e) quantum information processing, also provide (eb) a quantum advantage for (e) causal inference. Subjects: Quantum Physics (quant-ph); Learning (cs.LG); General Relativity and Quantum Cosmology (gr-qc); Machine Learning (stat.ML) Journal reference: Nat Phys 11, 414-420 (2015) DOI: 10.1038/nphys3266 Cite as: arXiv: 1406.5036 [quant-ph] (or arXiv: 1406.5036v1 [quant-ph] for this version) **Inferring causal structure: a quantum advantage Katja Ried, Megan Agnew, Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. ReschSTOP D. T. Son, A. O. Starinets** review recent progress in applying the AdS/CFT correspondence to (e&eb) finite-temperature field theory. In particular, they show how the hydrodynamic behavior of field theory is reflected in (e&eb) the low-momentum limit of (e) correlation functions computed through (e&eb) a real-time AdS/CFT prescription, which they formulate. Authors also show how the hydrodynamic modes in field theory correspond to (e&eb) the low-lying quasinormal modes of (e) the AdS black p-brane metric. They provide a proof of the universality of (e) the viscosity/entropy ratio within (eb) a class of theories with (e&eb) gravity duals and formulate (eb) a viscosity bound conjecture. Possible implications for real systems are mentioned. Subjects: High Energy Physics - Theory (hep-th) Journal reference: Ann.Rev.Nucl.Part.Sci.57:95-118,2007 DOI: 10.1146/annurev.nucl.57.090506.123120Report number: INT PUB 07-02Cite as: arXiv:0704.0240 [hep-th]

(or arXiv:0704.0240v2 [hep-th] for this version) An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society **Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP** Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv:

1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version **Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP Viscosity, Black Holes, and Quantum Field Theory D. T. Son, A. O. Starinets STOP Authors** study information retrieval from evaporating black holes, assuming that the internal dynamics of a black hole is unitary and rapidly mixing, and assuming that the retriever has unlimited control over the emitted Hawking radiation. If the evaporation of the black hole has already proceeded past the ``half-way" point, where (e) half of the initial entropy has been radiated away, then (eb) additional quantum information deposited in (eb) the black hole is revealed in (eb) the Hawking radiation very rapidly. Information deposited prior to the half-way point remains concealed until (e) the half-way point, and then emerges (eb) quickly. These conclusions hold because (e) typical local quantum circuits are efficient encoders for (e) quantum error-correcting codes that nearly achieve (eb) the capacity of the quantum erasure channel. Estimate of a black hole's information retention time, based on (e) speculative dynamical assumptions, is just barely compatible with (e&eb) the black hole complementarity hypothesis. **Black holes as mirrors: quantum information in random subsystems Patrick Hayden1 and John Preskill2 Published 26 September 2007 Journal of High Energy Physics, Volume 2007, JHEP09 (2007) STOP** By analyzing the key properties of black holes from the point of (e) view of quantum information, authors derive a model-independent picture of (e) black hole quantum computing. It has been noticed that this picture exhibits (eb) striking similarities with (e&eb) quantum critical condensates, allowing (eb) the use of a common language to describe (eb) quantum computing in both systems. Authors analyze such quantum computing by allowing coupling to (e&eb) external modes, under the condition that the external influence must be (=) soft-enough in order not to (e) offset the basic properties of the system. Authors derive model-independent bounds on some crucial time-scales, such as (=) the times of gate operation, decoherence, maximal entanglement and total scrambling. They show that for black hole type quantum computers all these time-scales are of (e) the order of the

black hole half-life time. Furthermore, authors construct explicitly a set of Hamiltonians that generates (eb) a universal set of quantum gates for (e) the black hole type computer. They find that the gates work at maximal energy efficiency. Furthermore, authors establish a fundamental bound on the complexity of quantum circuits encoded on (eb&e) these systems, and characterize the unitary operations that are (=) implementable. It becomes apparent that the computational power is very limited due to (e) the fact that the black hole life-time is of the same order of (e) the gate operation time. As a consequence, it is impossible to retrieve its information, within (eb) the life-time of a black hole, by externally coupling to (e&eb) the black hole qubits. However, authors show that, in principle, coupling to (e&eb) some of the internal degrees of freedom allows (eb) acquiring knowledge about the micro-state. Still, due to the trivial complexity of operations that can be performed, (eb) there is no time advantage over the collection of Hawking radiation and subsequent decoding. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Report number: LMU-ASC 03/16, MPP-2016-1 Cite as: arXiv:1605.01407 [hep-th] (or arXiv:1605.01407v1 [hep-th] for this version) **Universality of Black Hole Quantum Computing Gia Dvali, Cesar Gomez, Dieter Lust, Yasser Omar, and Benedikt Richter (Submitted on 4 May 2016) STOP** In this addendum to [arXiv: 1402.5674] two points are discussed. In the first additional evidence is provided (eb) for a dual connection between the geometric length of an Einstein-Rosen bridge and (e&eb) the computational complexity of the quantum state of the dual CFT's. The relation between growth of complexity and (e&eb) Page's “Extreme Cosmic Censorship” principle is also remarked on. The second point involves a gedanken experiment in which Alice measures (eb) a complete set of commuting observables at her end of (e) an Einstein-Rosen bridge is discussed. An apparent paradox is resolved by appealing to (e) the properties of GHZ tripartite entanglement. **Addendum to computational complexity and black hole horizons Leonard Susskind* Version of Record online: 18 JAN 2016 DOI: 10.1002/prop.201500093 STOP** Analyzing some well established facts, authors give a model-independent

parameterization of black hole quantum computing in terms of (e&*eb) a set of macro and micro quantities and their relations. These include the relations between the extraordinarily-small energy gap of (e) black hole qubits and important time-scales of (e) information-processing, such as, scrambling time and Page's time. Authors then show, confirming and extending previous results, that other systems of nature with (e&eb) identical quantum informatics features are (=) attractive Bose-Einstein systems at (eb) the critical point of quantum phase transition. Here authors establish a complete isomorphy between (e&eb) the quantum computational properties of these two systems. In particular, they show that the quantum hair of a critical condensate is (=) strikingly similar to (=) the quantum hair of a black hole. Irrespectively whether one takes the similarity between the two systems as a remarkable coincidence or as (=) a sign of a deeper underlying connection, the following is evident. Black holes are not unique in their way of (e) quantum information processing and they can manufacture black hole based quantum computers in labs by (e) taking advantage of quantum criticality. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Atomic Physics (physics. Atom-ph); Quantum Physics (quant-ph) Cite as: arXiv: 1601.01329 [hep-th] (or arXiv: 1601.01329v1 [hep-th] for this version) **Black Hole Based Quantum Computing in Labs and in the Sky Gia Dvali, Mischa Panchenko (Submitted on 6 Jan 2016) STOP Gia Dvali, Alexander Gußmann** show that the existence of black holes with (e&eb) classical skyrmion hair invalidates (e) standard proofs that global charges, such as (=) the baryon number, cannot be conserved by (e) a black hole. By carefully analyzing the standard arguments based on (e) a Gedankenexperiment in which a black hole is seemingly-unable to (e) return the baryon number that it swallowed, authors identify inconsistencies in this reasoning, which does not (e) take into the account neither the existence of skyrmion black holes nor the baryon/skyrmion correspondence. Authors then perform a refined Gedankenexperiment by incorporating (e) the new knowledge and show that (eb) no contradiction with conservation of baryon number takes place at (eb) any stage of black hole evolution. Analysis also indicates (eb) no conflict

between semi-classical black holes and (e&eb) the existence of baryonic gauge interaction arbitrarily-weaker than gravity. Next, authors study classical cross sections of a minimally-coupled massless probe scalar field scattered by (e) _ a skyrmion black hole. Authors investigate how the skyrmion hair manifests itself by (e) comparing this cross section with (e&eb) the analogous cross section caused by (e) a Schwarzschild black hole which has (e) the same ADM mass as (=) the skyrmion black hole. Here authors find an order-one difference in the positions of (e) the characteristic peaks in (eb) the cross sections. The peaks are shifted to smaller scattering angles when (e) the skyrmion hair is present. This comes from the fact that the skyrmion hair changes (e&eb) the near horizon geometry of the black hole when compared to (e&eb) a Schwarzschild black hole with same ADM mass. Authors keep the study of this second aspect general so that (e) the qualitative results obtained can also be applied to(e&eb) black holes with classical hair of different kind. Subjects: High Energy Physics - Theory (hep-th); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU-ASC 18/16 Cite as: arXiv:1605.00543 [hep-th] (or arXiv:1605.00543v2 [hep-th] for this version) A preliminary discussion is given of (e) the prospects that gravitational-wave observations of (e) binary inspiral of black holes could reveal or constrain (e&eb) quantum modifications to (e) black hole dynamics, such as are required to (e) preserve postulates of quantum mechanics. Different proposals for such modifications are characterized by (e) different scales, and the size of these scales relative to (e&eb) those probed by observation of (e) inspiral signals is important in determining (eb) the feasibility of finding experimental signatures. Certain scenarios with strong quantum modifications in (eb) a region extending well outside the horizon are expected to modify (e&eb) classical evolution, and distort (e&eb) the near-peak gravitational wave signal, suggesting (eb) a search for departures from (e) waveforms predicted by general relativity. The near agreement of the GW150914 signals with (e&eb) such waveforms is discussed, and indicates (eb) constraints on some such scenarios. Important strategies for more precise future tests are 1) to develop more precise

predictions from (e) scenarios proposing quantum modifications, and 2) searching for observed deviations from (e) numerical relativity predictions via (e&eb) analysis of gravity wave data, particularly focussing on (e&eb) the signal region corresponding to plunge and merger. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1602.03622 [gr-qc] (or arXiv: 1602.03622v2 [gr-qc] for this version) **Gravitational wave tests of quantum modifications to black hole structure -- with post-GW150914 update Steven B. Giddings STOP** Direct LHC bounds on (e&eb) colored SUSY particles now corner naturalness more than the measured value of the Higgs mass does. Bounds on the gluino are of (e) particular importance, since it radiatively "sucks" up (e) the stop and Higgs-up soft masses. As a result, even models that easily accommodate a 125 GeV Higgs are almost as tuned as (=) the simplest version of SUSY, the MSSM: at best at the percent level. In this paper, authors further examine how current LHC results constrain (e) naturalness in three classes of models that may relax LHC bounds on (e&eb) sparticles: split families (e&eb), baryonic RPV (e&eb), and Dirac gauginos. In models of split families and bRPV, the bounds on the gluino are (=) only slightly reduced, resulting in (eb) a few percent tuning. In particular, having a natural spectrum in bRPV models typically implies (eb) that tops, Ws, and Zs are easily produced (eb) in the cascade decays of (e) squarks and gluinos. The resulting leptons and missing energy push (eb) the gluino mass limit above 1 TeV. Even when the gluino has a Dirac mass and does not contribute to the stop mass at one loop, tuning reappears in calculable models because there is no symmetry imposing the supersoft limit. Authors conclude that, even if sparticles are found at LHC-14, naturalness will not emerge triumphant. Subjects: High Energy Physics - Phenomenology (hep-ph)DOI:10.1007/JHEP03(2014)022 Cite as:rXiv:1309.3568 [hep-ph] (or arXiv:1309.3568v2 [hep-ph] for this version) **The Last Vestiges of Naturalness Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang, Ken Van Tilburg, Giovanni Villadoro STOP** The lack of evidence for new physics beyond (e) the standard model at the LHC points to (e) a paucity of new particles near the weak scale. This suggests that the weak scale is tuned and that supersymmetry, if present at

all, is realized (eb) at higher energies. The measured Higgs mass constrains (e) the scalar sparticles to be below 10^5 TeV, while (e) gauge coupling unification favors (eb) Higgsinos below 100 TeV. Nevertheless, in many models gaugino masses are (=) suppressed and remain within reach of the LHC. Tuning the weak scale and (e&eb) the renormalization group evolution of (e) the scalar masses constrain (e) Split model building. Due to the small gaugino masses, either the squarks or the up-higgs often run tachyonic; in the latter case, successful electroweak breaking requires heavy higgsinos near the scalar sparticles. We discuss the consequences of tuning the weak scale and the phenomenology of several models of Split supersymmetry including anomaly mediation, U(1)_(B-L) mediation, and Split gauge mediation. Subjects: High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Experiment (hep-ex) DOI: 10.1007/JHEP02 (2013)126 Cite as: arXiv: 1210.0555 [hep-ph] (or arXiv: 1210.0555v2 [hep-ph] for this version) **Mini-Split Asimina Arvanitaki, Nathaniel Craig, Savas Dimopoulos, Giovanni Villadoro STOP** Light states associated with (e&eb) the hierarchy problem affect (e&eb) the Higgs LHC production and decays. Authors illustrate this within the MSSM and two simple extensions applying (e&eb) the latest bounds from (e) LHC Higgs searches. Large deviations in the Higgs properties are expected in (eb) a natural SUSY spectrum. The discovery of a non-Standard-Model Higgs may signal (eb) the presence of light stops accessible at (eb) the LHC. Conversely, the more the Higgs is Standard-Model-like, (eb) the more tuned the theory becomes. Taking the ratio of different Higgs decay channels at the LHC cancels (er&eb) the leading (eb) QCD uncertainties and potentially improves (eb) the accuracy in Higgs coupling measurements to (e) the percent level. This may lead to (eb) the possibility of doing precision Higgs physics at (eb) the LHC. Finally, authors entertain the possibility that the ATLAS excess around (e&eb) 125 GeV persists with (e&eb) a Higgs production cross-section that is enhanced compared to (e&eb) the SM. This increase can only be accommodated in extensions of (e) the MSSM and it may suggest (eb) that stops lie below 400 GeV, likely within (eb) reach of next year's LHC run. Subjects: High Energy Physics - Phenomenology (hep-ph) DOI: 10.1007/JHEP02 (2012)144 Cite

as: arXiv: 1112.4835 [hep-ph] (or arXiv: 1112.4835v2 [hep-ph] for this version) **A Non Standard Model Higgs at the LHC as a Sign of Naturalness Asimina Arvanitaki, Giovanni Villadoro STOP** It has recently been suggested that the presence of a plenitude of light axions, an Axiverse, is (=) evidence for the extra dimensions of string theory. Authors discuss the observational consequences of (e) these axions on astrophysical black holes through (e&eb) the Penrose superradiance process. When an axion Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e&eb) the black hole "nucleus" forming (eb) a gravitational atom in the sky. The occupation number of superradiant atomic levels fed by (e) the energy and angular momentum of the black hole grows exponentially. The black hole spins down and an axion Bose-Einstein condensate cloud forms (eb) around it. When the attractive axion self-interactions become (=) stronger than the gravitational binding energy, (eb) the axion cloud collapses, a phenomenon known in condensed matter physics as "Bosenova". The existence of axions is first diagnosed by (e) gaps in the mass vs (e&eb) spin plot of astrophysical black holes. For young black holes the allowed values of (e) spin are quantized, giving rise to (eb) "Regge trajectories" inside the gap region. The axion cloud can also be observed directly either through (e&eb) precision mapping of the near horizon geometry or through (e&eb) gravitational waves coming from (e) the Bosenova explosion, as well as axion transitions and (e&eb) annihilations in (eb) the gravitational atom. Estimates suggest (eb) that these signals are detectable in (eb) upcoming experiments, such as Advanced LIGO, AGIS, and LISA. Current black hole spin measurements imply (eb) an upper bound on (e&eb) the QCD axion decay constant of (e) 2 x 10^17 GeV, while (e) Advanced LIGO can detect (e&eb) signals from (e) a QCD axion cloud with (e&eb) a decay constant as low as (=) the GUT scale. Authors finally discuss the possibility of observing the gamma-rays associated with the Bosenova explosion and, perhaps, the radio waves from axion-to-photon conversion for the QCD axion. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Journal

reference:Phys.Rev.D83:044026, 2011 DOI: 10.1103/PhysRevD.83.044026 Cite as: arXiv: 1004.3558 [hep-th] (or arXiv: 1004.3558v2 [hep-th] for this version) **Exploring the String Axiverse with Precision Black Hole Physics Asimina Arvanitaki, Sergei Dubovsky (Submitted on 20 Apr 2010 (v1), last revised 16 Aug 2011 (this version, v2)) STOP Skyrmion Black Hole Hair: Conservation of Baryon Number by Black Holes and Observable Manifestations Gia Dvali, Alexander Gußmann STOP Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust** discuss BMS supertranslations both at (eb) null-infinity and on the horizon for the case of the Schwarzschild black hole. Authors show that both kinds of supertranslations lead to (eb) infinitely many gapless physical excitations. On this basis authors construct a quotient algebra using (e) suited superpositions of (e) both kinds of transformations which cannot be compensated by (e) an ordinary BMS-supertranslation and therefore are (=) intrinsically due to (e) the presence of an event horizon. Authors show that these quotient transformations are (=) physical and generate (eb) gapless excitations on the horizon that can account for (e) the gravitational hair as well as for (e) the black hole entropy. They identify the physics of these modes as associated with (e&eb) Bogolioubov-Goldstone modes due to quantum criticality. Classically the number of these gapless modes is (=) infinite. However, authors show that due to (e) quantum criticality the actual amount of information-carriers becomes (=) finite and consistent with Bekenstein entropy. Although they only consider the case of Schwarzschild geometry, the arguments are extendable to(e&eb) arbitrary space-times containing event horizons. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU--ASC 05/16, MPP--2016—3 Cite as: arXiv: 1601.03725 [hep-th] (or arXiv: 1601.03725v1 [hep-th] for this version) **Gravitational Black Hole Hair from Event Horizon Supertranslations Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust STOP Gia Dvali, Cesar Gomez, Nico Wintergerst** suggest that holography can be formulated in terms of (e&eb) the information capacity of (e) the St\"uckelberg degrees of freedom that maintain

gauge invariance of (e) the theory in the presence of (e) an information boundary. These St\"uckelbergs act as (e&eb) qubits that account for a certain fraction of (e) quantum information. Their information capacity is measured by (e) the ratio of the inverse St\"uckelberg energy gap to (e) the size of the system. Systems with (e&eb) the smallest gap are (=) maximally holographic. For massless gauge systems this information measure is (=) universally equal to (=) the inverse coupling evaluated at the systems' length scale. In this language it becomes very transparent why the St\"uckelberg information capacity of black holes saturates (e full) the Bekenstein bound and accounts for (e) the entire information of the system. The physical reason is that the strength of quantum interaction is bounded from below by the gravitational coupling, which scales as (=) area. Observing the striking similarity between the scalings of the energy gap of the boundary St\"uckelberg modes and (e&eb): (AC=DC=0) the Bogoliubov modes of critical many-body systems, authors establish a connection between holography and (e&eb) quantum criticality through (e&eb) the correspondence between these modes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1511.03525 [hep-th] (or arXiv: 1511.03525v1 [hep-th] for this version) **Stückelberg Formulation of Holography Gia Dvali, Cesar Gomez, Nico Wintergerst STOP B**lack hole entropy, denoted by N, in (semi)classical limit is (=) infinite. This scaling reveals (eb) a very important information about the qubit degrees of (e) freedom that carry black hole entropy. Namely, the multiplicity of qubits scales as N, whereas (e) their energy gap and their coupling as 1/N. Such a behavior is indeed exhibited by (e) Bogoliubov-Goldstone degrees of freedom of (e) a quantum-critical state of N soft gravitons (a condensate or a coherent state) describing (eb) the black hole quantum portrait. They can be viewed as (=) the Goldstone modes of (e) a broken symmetry acting on (e&eb) the graviton condensate. In this picture Minkowski space naturally emerges as (=) a coherent state of infinite-N gravitons of (e) infinite wavelength and it carries (e&eb) an infinite entropy. In this paper authors **Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015)** ask what is the geometric meaning (if any) of (e) the classical limit of this symmetry. Authors argue that the infinite-

N limit of Bogoliubov-Goldstone modes of critical graviton condensate is described by (e) recently-discussed classical BMS super-translations broken by the black hole geometry. However, the full black hole information can only be recovered for (e) finite N, since the recovery time becomes (=) infinite in classical limit in which N is (=) infinite. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) DOI: 10.1016/j.physletb.2015.11.073 Report number: MPP-2015-215, LMU-ASC 59/15 Cite as: arXiv: 1509.02114 [hep-th] (or arXiv: 1509.02114v1 [hep-th] for this version) **Classical Limit of Black Hole Quantum N-Portrait and BMS Symmetry Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015) STOP** Recent ideas about understanding physics of black hole information-processing in terms of (e&eb) quantum criticality allow us to (eb) implement black hole mechanisms of quantum computing within (eb) critical Bose-Einstein systems. The generic feature, uncovered both by analytic and numeric studies, is (=) the emergence at the critical point of (e) gapless weakly-interacting modes, which act as (=) qubits for information-storage at a very low energy cost. These modes can be effectively described in terms of (e&eb) either Bogoliubov or Goldstone degrees of freedom. The ground-state at the critical point is (=) maximally entangled and far from being classical. We confirm this near-critical behavior by a new analytic method. **Gia Dvali, Mischa Panchenko** compute growth of entanglement and show its consistency with (e&eb) black hole type behavior. On the other hand, in the over-critical regime the system develops (eb) a Lyapunov exponent and scrambles quantum information very fast. By, manipulating (e&eb) the system parameters externally, authors can put it in and out of various regimes and in this way control (e&eb) the sequence of information storage and processing. By using gapless Bogoliubov modes as (=) control qubits they design some simple logic gates. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Cite as: arXiv: 1507.08952 [hep-th] (or arXiv: 1507.08952v1 [hep-th] for this version) **Black Hole Type Quantum**

**Computing in Critical Bose-Einstein Systems Gia Dvali, Mischa Panchenko (Submitted on 31 Jul 2015) STOP Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst** establish a fundamental connection between quantum criticality of a many-body system, such as Bose-Einstein condensates, and (e&eb) its capacity of information-storage and processing. For deriving the effective theory of modes in the vicinity of the quantum critical point they develop a new method by mapping a Bose-Einstein condensate of N-particles onto (e&eb) a sigma model with a continuous global (pseudo)symmetry that mixes (e&eb) bosons of different momenta. The Bogolyubov modes of the condensate are mapped onto (e&eb) the Goldstone modes of the sigma model, which become (=) gapless at the critical point. These gapless Goldstone modes are (=) the quantum carriers of information and entropy. Analyzing their effective theory, authors observe the information-processing properties strikingly similar to (=) the ones predicted by the black hole portrait. The energy cost per qubit of information-storage vanishes in (eb) the large-N limit and the total information-storage capacity increases with (e&eb) N either exponentially or as a power law. The longevity of information-storage also increases with (e&eb) N, whereas (e) the scrambling time in (eb) the over-critical regime is controlled by (e&eb) the Lyapunov exponent and scales logarithmically with (e&eb) N. This connection reveals (eb) that the origin of black hole information storage lies in (eb) the quantum criticality of the graviton Bose-gas, and that much simpler systems that can be manufactured in table-top experiments can exhibit (eb) very similar information-processing dynamics. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 92, 125002 (2015) DOI: 10.1103/PhysRevD.92.125002 Cite as: arXiv:1507.02948 [hep-th] (or arXiv:1507.02948v1 [hep-th] for this version) **Nambu-Goldstone Effective Theory of Information at Quantum Criticality Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst STOP Gia Dvali, Cesar Gomez** show that a positive cosmological constant is (=) incompatible with the quantum-corpuscular resolution of de Sitter metric in form of (e) a coherent state. The reason is very general and is due to (e) the

quantum self-destruction of (e) the coherent state because of (e) the scattering of constituent graviton quanta. This process creates an irreversible quantum clock, which precludes (e) eternal de Sitter. It also eliminates (e) the possibility of Boltzmann brains and Poincare recurrences. This effect is expected to be (=) part of any microscopic theory that takes into account (e) the quantum corpuscular structure of the cosmological background. This observation puts the cosmological constant problem in a very different light, promoting it, from a naturalness problem, into (e&eb) a question of quantum consistency. We are learning that quantum gravity cannot tolerate (e) exceedingly-classical sources. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1412.8077 [hep-th]

(or arXiv:1412.8077v1 [hep-th] for this version) **Quantum Exclusion of Positive Cosmological Constant? Gia Dvali, Cesar Gomez (Submitted on 27 Dec 2014) STOP G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger** establish a connection between the ultra-Planckian scattering amplitudes in field and string theory and (e&eb) unitarization by black hole formation in these scattering processes. **Using as a guideline an explicit microscopic theory in which the black hole represents (eb) a bound-state of many soft gravitons at (eb) the quantum critical point, they were able to identify and compute (e&eb) a set of perturbative amplitudes relevant for (e) black hole formation. These are the tree-level N-graviton scattering S-matrix elements in (eb) a kinematical regime (called classicalization limit) where (e) the two incoming ultra-Planckian gravitons produce (eb) a large number N of soft gravitons. Authors compute these amplitudes by using (e) the Kawai-Lewellen-Tye relations, as well as scattering equations and (e&eb) string theory techniques. Authors discover that this limit reveals (eb) the key features of the microscopic corpuscular black hole N-portrait. In particular, the perturbative suppression factor of a N-graviton final state, derived from (e) the amplitude, matches the non-perturbative black hole entropy when (e) N reaches the quantum criticality value, whereas final states with different value of N are (=) either suppressed or excluded by non-perturbative**

**corpuscular physics. Thus they identify the microscopic reason behind the black hole dominance over other final states including (e) non-black hole classical object. In the parameterization of the classicalization limit the scattering equations can be solved exactly allowing (eb) us to obtain closed expressions for the high-energy limit of the open and closed superstring tree-level scattering amplitudes for a generic number N of external legs. Authors demonstrate matching and complementarity between the string theory and (e&eb) field theory** in different large-s and large-N regimes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1016/j.nuclphysb.2015.02.004 Report number: MPP-2014-320, LMU-ASC 52/14 Cite as: arXiv:1409.7405 [hep-th] (or arXiv:1409.7405v2 [hep-th] for this version) **Black Hole Formation and Classicalization in Ultra-Planckian 2 -> N Scattering G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger STOP** Gravitational backgrounds, such as black holes, AdS, de Sitter and inflationary universes, should be viewed as (=) composite of N soft constituent gravitons. It then follows that such systems are close to (e) quantum criticality of graviton Bose-gas to Bose-liquid transition. Generic properties of the ordinary metric description, including (e) geodesic motion or particle-creation in the background metric, emerge as (=) the large-N limit of quantum scattering of (e) constituent longitudinal gravitons. **Gia Dvali, Cesar Gomez** show that this picture correctly accounts for (e) physics of large and small black holes in AdS, as well as reproduces (eb) well-known inflationary predictions for cosmological parameters. However, it anticipates new effects not captured by (e) the standard semi-classical treatment. In particular, **Gia Dvali, Cesar Gomez** predict observable corrections that are sensitive to (e&eb) the inflationary history way beyond last 60 e-foldings. They derive an absolute upper bound on the number of e-foldings, beyond (e) which neither de Sitter nor inflationary Universe can be approximated by (e) a semi-classical metric. However, they could in principle persist in a new type of (e) it quantum eternity state. Authors discuss implications of this phenomenon for (e&eb) the cosmological constant problem. Subjects: High Energy Physics - Theory (hep-th); Cosmology and

Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Quantum Physics (quant-ph) DOI: 10.1088/1475-7516/2014/01/023 Cite as: arXiv: 1312.4795 [hep-th] (or arXiv:1312.4795v1 [hep-th] for this version) **Quantum Compositeness of Gravity: Black Holes, AdS and Inflation Gia Dvali, Cesar Gomez STOP** As quantum information science approaches the goal of constructing quantum computers, understanding loss of information through (e&eb) decoherence becomes (=) increasingly important. The information about a system that can be obtained from (e) its environment can facilitate quantum control and error correction. Moreover, observers gain (eb) most of their information indirectly, by (e) monitoring (primarily photon) environments of (e) the "objects of interest." Exactly how this information is inscribed in the environment is essential for (e) the emergence of "the classical" from (e) the quantum substrate. In this paper, **Robin Blume-Kohout, Wojciech H. Zurek** examine how many-qubit (or many-spin) environments can store (e) information about a single system. The information lost to (e) the environment can be stored redundantly, or it can be encoded in (eb) entangled modes of the environment. Authors go on to show that randomly chosen states of the environment almost always encode (e) the information so that (e) an observer must capture (e) a majority of the environment to deduce (eb) the system's state. Conversely, in the states produced by (e) a typical decoherence process, information about a particular observable of the system is stored redundantly. This selective proliferation of "the fittest information" (known as Quantum Darwinism) plays a key role in (eb) choosing the preferred, effectively classical observables of (e) macroscopic systems. The developing appreciation that the environment functions not just as (=) a garbage dump, but as a communication channel, is extending understanding of the environment's role in (e&eb) the quantum-classical transition beyond (e) the traditional paradigm of decoherence. Keywords Decoherence quantum Darwinism quantum information quantum-classical transition **Foundations of Physics November 2005, Volume 35, Issue 11, pp 1857-1876 A Simple Example of "Quantum Darwinism": Redundant Information Storage in Many-Spin Environments Robin Blume-Kohout, Wojciech H. Zurek STOP Harold Ollivier, David**

**Poulin, and Wojciech H. Zurek** study the emergence of objective properties in open quantum systems. In our analysis, the environment is promoted from a passive role of a reservoir selectively destroying quantum coherence to an active role of amplifier selectively **proliferating information about the system. We show that only preferred pointer states of the system can leave a redundant and therefore easily detectable imprint on the environment. Observers who—as is almost always the case—discover the state of the system indirectly (by probing a fraction of its environment) will find out only about the corresponding pointer observable. Many observers can act in this fashion independently and without perturbing the system. They will agree about its state. In this operational sense, preferred pointer states exist objectively. Received 24 July 2003 DOI:http://dx.doi.org/10.1103/PhysRevLett.93.220401 Objective Properties from Subjective Quantum States: Environment as a Witness Harold Ollivier, David Poulin, and Wojciech H. Zurek Phys. Rev. Lett. 93, 220401 – Published 22 November 2004 STOP** Quantum Darwinism recognizes that we—the observers—acquire our information about the "systems of interest" indirectly from their imprints on the environment. **Michael Zwolak, H. T. Quan, and Wojciech H. Zurek** show that information about a system can be acquired from a mixed-state, or hazy, environment, but the storage capacity of an environment fragment is suppressed by its initial entropy. In the case of good decoherence, the mutual information between the system and the fragment is given solely by the fragment's entropy increase. For fairly mixed environments, this means a reduction by a factor 1−h, where h is the haziness of the environment, i.e., the initial entropy of an environment qubit. Thus, even such hazy environments eventually reveal the state of the system, although now the intercepted environment fragment must be larger by ~ (1−h)−1 to gain the same information about the system. DOI:http://dx.doi.org/10.1103/PhysRevLett.103.110402 **Quantum Darwinism in a Mixed Environment Michael Zwolak, H. T. Quan, and Wojciech H. Zurek Phys. Rev. Lett. 103, 110402 – Published 8 September 2009** We study quantum Darwinism—the redundant recording of information about the preferred states of a decohering system by its environment—for an object illuminated by a

blackbody. In the cases of point-source and isotropic illumination, we calculate the quantum mutual information between the object and its photon environment. We demonstrate that this realistic model exhibits fast and extensive proliferation of information about the object into the environment and results in redundancies orders of magnitude larger than the exactly soluble models considered to date. DOI:http://dx.doi.org/10.1103/PhysRevLett.105.020404 **Quantum Darwinism in an Everyday Environment: Huge Redundancy in Scattered Photons C. Jess Riedel and Wojciech H. Zurek Phys. Rev. Lett. 105, 020404 – Published 9 July 2010 STOP** Measurements transfer information about a system to the apparatus and then, further on, to observers and (often inadvertently) to the environment. I show that even imperfect copying essential in such situations restricts possible unperturbed outcomes to an orthogonal subset of all possible states of the system, thus breaking the unitary symmetry of its Hilbert space implied by the quantum superposition principle. Preferred outcome states emerge as a result. They provide a framework for "wave-packet collapse," designating terminal points of quantum jumps and defining the measured observable by specifying its eigenstates. In quantum Darwinism, they are the progenitors of multiple copies spread throughout the environment—the fittest quantum states that not only survive decoherence, but subvert the environment into carrying information about them—into becoming a witness. DOI:http://dx.doi.org/10.1103/PhysRevA.76.052110 **Quantum origin of quantum jumps: Breaking of unitary symmetry induced by information transfer in the transition from quantum to classical Wojciech Hubert Zurek Phys. Rev. A 76, 052110 – Published 16 November 2007 STOP** The Born rule, a foundational axiom used to deduce probabilities of events from **wavefunctions, is indispensable in the everyday practice of quantum physics. It is also key in the quest to reconcile the ostensibly inconsistent laws of the quantum and classical realms, as it confers physical significance to reduced density matrices, the essential tools of decoherence theory. Following Bohr's Copenhagen interpretation, textbooks postulate the Born rule outright. However, recent attempts to derive it from (e) other quantum principles have been successful, holding promise for (e)**

**simplifying and clarifying the quantum foundational bedrock. A major family of derivations is based on envariance, a recently discovered symmetry of (e) entangled quantum states. Here, authors identify and experimentally test three premises central to (e) these envariance-based derivations, thus demonstrating, (eb) in the microworld, the symmetries from (e) which the Born rule is derived. Further, they demonstrate envariance in a purely local quantum system, showing (eb) its independence from relativistic causality.Subjects: Quantum Physics (quant-ph)Cite as:arXiv:1604.01471 [quant-ph] (or arXiv:1604.01471v1 [quant-ph] for this version) Quantum probabilities from quantum entanglement: Experimentally unpacking the Born rule Jérémie Harris, Frédéric Bouchard, Enrico Santamato, Wojciech H. Zurek, Robert W. Boyd, Ebrahim Karimi**

Integrating $[m_4]_4(t, a_4)$ with respect to a from 0 to ∞, we have

$$\frac{dM_4(t)}{dt} = \int_0^\infty [mt]_4(t, a_4)da_4$$

$$= -\int_0^\infty [m_a]_4\,(t, a_4)da_4 - \int_0^\infty \mu(a_4)m(t, a_4)da_4$$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f^{'}{}_{4*}]_4 \int_0^\infty p_4^*(a_4)da_4\right)M_4 - [f_{2*}^4]_4 \int_0^\infty p_4^*(a_4)da_4 W_4$$

$= m_4(t, 0) - \int_0^\infty \mu_4(a_4)\, m_4(t, a_4)da_4$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}^{i}]_4 P_4^*\right)M_4 - [f_{2*}^{i}]_4 P_4^* W_{\cdot 4}$$

From (30), we arrive at the equations

$$\frac{dW_4(t)}{dt} = \int_0^{\infty} ([h_*]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4$$

$$-\left([f_{4*}] + [f_{2*}]_4 + \left[f_{4*}^{'}\right]_4 P_4^*\right) M_4 + \left(\frac{\left[h_*^{'}\right]_4}{[h_*]^4}\, P_4^*(0) - f_{2*}^{'} P_4^*\right) W_4,$$

(34)

$$\frac{dW_4(t)}{dt}\left[g_*^{'}\right]_4 Q_4^*\, M_{.4}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{r*}]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4 = ([h_*]_4\beta_4(\theta_4) - \mu_4(\theta_4))M_{,4}$$

Where $0< \theta_4 = \theta_4(t) < \infty$.

And if $a_4^2\,[a_4^2]_4<4\,[a_2]_4<1\,a_2 e$,

$\mathrm{Re}\,\lambda_4\left[\lambda_4\right]_4 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_4 = -\frac{4}{2}a_4[a2\ ]_4 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_*\,[h_*]_4\beta_4(a_4) - \mu_4(a_4)\right|$. Then$|n_4(t)|$ ⊠⊿∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_4^2 <4\, a_2 e$ and $[a_4] > 2\Delta$,or $[a_4^2]$ ⊠⊿$[a_2]e$ and $(a_4 - \sqrt{(a_4^2)_4 - 4(a_2)_4 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+4}\right)$ and $\beta_4(a_4) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4) + d\tau_4 da} = \frac{4}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_4 + [\mu_0]$ 0.5. Assume$[f_4](x_4)= K_4 x_4$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_4(x_4)= [0x_4]/(1 + x_4)$; the functional responses $[[f_2](x) = [0]\, x_4/(1 + x_4)$and $h_4(x_4)=(1+\omega_4 x_4)/(1 + x_4)$(See Holling [7] and Rashevsky [13])

Equation (44),now has the form of

$$\frac{m_4 + [\mu_0] + [K][P_0])}{[m_4]_4 + (\mu_0)_4 + [K][P_0])\,{}_4^3\,[K]_4[P_0]_4^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(4 + P_0) < b_4$ i e $[P_0] < b/(\delta_4 - b_4)$then the equilibrium$(P_0)$ $[P_0]_4(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_4$

$[f_{4*}] = K_4\, P^*\, [f_{2*}] = \frac{\sigma Q_4^*}{4+Q_4^*} \quad [h_*] = \frac{4+\omega_4 Q_4^*}{4+Q_4^*}$

$[f_{4*}] = K_4\, [f_{2*}^4] = \frac{\sigma_4}{(4+Q_4^*)^2}\, [g_*^4] = \frac{\sigma_4}{(4+P_4^*)^2} \quad [h_*^4] = \frac{\omega_4 - 4}{(4+Q_4^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0$, K, b, $\delta$, and $\omega$ need to satisfy

$$\frac{\omega_4}{(m + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma)^2} + \frac{2\omega_4}{([m_4]_4 + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma_4)^3} < 1$$

And

$$\frac{1}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^2}+\frac{2\omega_4}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4}+(m_4+[\mu_0]_4+\frac{K_4b_4}{\delta_4-b_4}+\sigma_4)\,Q_4^*)^3$$

$$=m_4+[\mu_0]_4+3+\frac{K_4b_4}{\delta_4-b_4}(4+Q^*)(1+\omega_4Q_4^*$$

Now

$$\left[a_4^2\right]-1\left[a_2\right]e=\left(\frac{2K_4b_4}{\delta_4-b_4}\right)^2+\frac{1K_4b_4\sigma_4\,Q_4^*}{(\delta_4-b_4)(1+Q_4^*)}+\frac{\sigma_4^2Q_4^{*2}}{(1+Q_4^*)^2}$$

$$\frac{1\,b_4\sigma_4\,(\delta_4-b)Q^*}{\delta_4(1+Q_4^*)^2}-\frac{1\,(\delta_4-b_4)^2(1-\omega_4)Q_4^*}{\delta_4(1+\omega_4Q_4^*)(1+Q_4^*)}P^*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_4\,ae_4^{mu}\;[\mu_0]_4-\frac{a_4}{a_4+4}$$

Hence if $[a_4^2>4(a_2)_4e$ and

$$[a_4] = \frac{2K_4 b_4}{\delta_4 - b_4} + \frac{\sigma_4 \, Q_4^*}{1 + Q_4^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_4]_4 = [f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}] P_4^*$$

$$[a_2]_4 = \left[f_{2*}^4\right] P_4^* - \frac{h_*^4}{h_*} P^*(0)$$

$$e = [g_*^4]_4 \, Q_4^*$$

$$\eta_4(t) = \left[h_*\right]_4 \beta_4 \left(\theta_4(t)\right) - \mu_4(\theta(t)),$$

The system(34)can be written as the following system:

$$[M(t)]\left[a_4\right] M_4 - [a_2]_4 W_4 - [n(t)] M_4$$

$$w_4(t) = eM_4$$

Write

$$A = \begin{pmatrix} -[a_4]_4 - [a_2]_4 & \\ e & 0 \end{pmatrix}$$

$$B = (t) = \begin{pmatrix} [\eta(t)]_4 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_4 \\ W_4 \end{pmatrix}$$

The solution of (32)is

$$X_4(t) = e^{A,t}[X]_4(0) + \int_0 e^{At}$$

**Some of the other conceptual and thematic and discursive factorial representationalities to which the model is applicable are:every one-dimensional quantum mechanical Hamiltonian H1 can have (e) a partner H2 such that H1 and H2 taken together may be viewed as (=) the components of a supersymmetric Hamiltonian. The term 'supersymmetric Hamiltonian' is taken to mean (eb) a Hamiltonian defined in terms of (e&eb)charges that obey (eb)the same algebra as that of (e) the generators of supersymmetry in field theory. The consequences of this symmetry for (er) the spectra of H1 and H2 are explored. It is shown how the supersymmetric pairing may be utilised to eliminate (e) the ground state of H1, or add a state below the ground state of H1 or maintain the spectrum of H1. It is also explicitly demonstrated that the supersymmetric pairing may be used to generate (eb) a class of anharmonic potentials with (e&eb) exactly specified spectra. The complete spectrum of an anharmonic potential so generated consists of (e) all the eigenstates of the simple harmonic oscillator and, in addition, a ground state at a specified energy E which lies arbitrarily below (eb) the E=1/2 ground state of the harmonic oscillator. Supersymmetric quantum mechanics of one-dimensional systems C V Sukumar Journal of Physics A: Mathematical and General, Volume 18, Number 15 STOP Authors examine the properties and consequences of (e&eb) pseudo-supersymmetry for (e) quantum systems admitting (e) a pseudo-Hermitian Hamiltonian. Authors explore the Witten index of (e) pseudo-**

**supersymmetry and show that every pair of diagonalizable (not necessarily Hermitian) Hamiltonians with (e&eb) discrete spectra and (e&eb) real or complex-conjugate pairs of (e) eigenvalues are (=) isospectral and have (e) identical degeneracy structure except (e) perhaps for the zero eigenvalue if and only if (e) they are pseudo-supersymmetric partners. This implies (eb) that pseudo-supersymmetry is (=) the basic framework for (e) generating non-Hermitian PT-symmetric and (e&eb) non-PT-symmetric Hamiltonians with (e&eb) a real spectrum via (e&eb) a Darboux transformation, and shows (eb) that every diagonalizable Hamiltonian H with(e&eb) a discrete spectrum and (e&eb) real or complex-conjugate pairs of (e) eigenvalues may be factored as (=) H=L♯L where L is (=) a linear operator with pseudo-adjoint L♯. In particular, this factorization applies to (e&eb) PT-symmetric and Hermitian Hamiltonians. The non-degenerate two-level systems provide (eb) a class of Hamiltonians that are pseudo-Hermitian. Authors demonstrate the implications of general results for (e&eb) this class in some detail. Copyright © 2002 Elsevier Science B.V. All rights reserved. Nuclear Physics B Volume 640, Issue 3, 30 September 2002, Pages 419–434Pseudo-supersymmetric quantum mechanics and isospectral pseudo-Hermitian Hamiltonians Ali Mostafazadeh STOP Michele Caselle, Alessandro Nada, Marco Panero present a high-precision lattice calculation of (e) the equation of state in the confining phase of (e) SU(2) Yang-Mills theory. Authors show that the results are described very well by (e) a gas of massive, non-interacting glueballs, provided (e) one assumes an exponentially growing Hagedorn spectrum. The latter can be derived within (eb) an effective bosonic closed-string model, leading to (eb) a parameter-free theoretical prediction, which is in perfect agreement with (=) proposed lattice results. Furthermore, when applied to (e&eb) SU (3) Yang-Mills theory, this effective model accurately describes (eb) the lattice results reported by Bors\'anyi et al. in JHEP 07 (2012) 056. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 07 (2015) 143 DOI: 10.1007/JHEP07 (2015)143 Cite as: arXiv:**

**1505.01106 [hep-lat] (or arXiv: 1505.01106v2 [hep-lat] for this version) Hagedorn spectrum and thermodynamics of SU (2) and SU (3) Yang-Mills theories Michele Caselle, Alessandro Nada, Marco Panero STOP B. de Wit J. Hoppe, H. Nicolai study the quantum-mechanical properties of (e) a supermembrane and examine the nature of (e) its ground state. A supersymmetric gauge theory of (e) area-preserving transformations provides (eb) a convenient framework for this study. The supermembrane can be viewed as (=) a limiting case of a class of models in supersymmetric quantum mechanics. Its mass does not depend on (=) the zero modes and vanishes only if (e) the wave function is (=) a singlet under (e&eb) supersymmetry transformations of (e) the nonzero modes. Authors exhibit (eb) the complexity of (e) the supermembrane ground state and (e&eb) examine various truncations of these models. None of these truncations has massless states. Copyright © 1988 Published by Elsevier B.V. Nuclear Physics B Volume 305, Issue 4, 28 December 1988: Pages 545-581: On the quantum mechanics of supermembranes Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. B. de Wit J. Hoppe, H. Nicolai STOP In this note we explain how world-volume geometries of (e) D-branes can be reconstructed within (eb) the microscopic framework where (e) D-branes are described through (e&eb) boundary conformal field theory. Authors extract the (non-commutative) world-volume algebras from (e) the operator product expansions of (e) open string vertex operators. For branes in a flat background with (e&eb) constant non-vanishing B-field, the operator products are computed perturbatively to (e) all orders in the field strength. Resulting series coincides with (=) Kontsevich's presentation of the Moyal product. After extending these considerations to (e) fermionic fields authors conclude with some remarks on the generalization of approach to (e&eb) curved backgrounds. D-branes and deformation quantization Volker Schomerus1,2 Published 25 July 1999 Journal of High Energy Physics, Volume 1999, JHEP06(1999)STOP Authors employ the path integral approach developed in Gangopadhyay S. and Scholtz F. E., Phys. Rev.**

**Lett., 102 (2009) 241602 to discuss the (generalized) harmonic oscillator in (eb) a noncommutative plane. The action for (e) this system is derived in (e&eb) the coherent state basis with (e&eb) additional degrees of freedom. From this the action in the coherent state basis without (e) any additional degrees of freedom is obtained. This gives (eb) the ground-state spectrum of the system. Authors then employ the exact renormalization group approach to show (eb) that equivalence can be constructed between this (noncommutative) system and (e&eb) a commutative system Path integral action and exact renormalization group dualities for quantum systems in noncommutative plane Sunandan Gangopadhyay Published 25 June 2015 • Copyright © EPLA, 2015 STOP Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini present a lattice study of (e)the equation of state in Yang-Mills theory based on (e) the exceptional G(2) gauge group. As is well-known, at zero temperature this theory shares (e&eb) many qualitative features with real-world QCD, including (e) the absence of colored states in the spectrum and dynamical string breaking at (e) large distances. In agreement with previous works, authors show that at finite temperature this theory features (eb) a first-order deconfining phase transition, whose nature can be studied by (e) a semi-classical computation. Authors also show that the equilibrium thermodynamic observables in the deconfined phase bear striking quantitative similarities with (=) those found in SU (N) gauge theories: in particular, these quantities exhibit (eb) nearly perfect proportionality to the number of gluon degrees of freedom, and (e&eb) the trace anomaly reveals (eb) a characteristic quadratic dependence on (e&eb) the temperature, also observed in SU (N) Yang-Mills theories (both in four and in three spacetime dimensions). Authors compare our lattice data with (e&eb) analytical predictions from (e) effective models, and discuss their implications for (e) the deconfinement mechanism and (e&eb) high-temperature properties of (e) strongly interacting, non-supersymmetric gauge theories. Results give (eb) strong evidence for (e) the conjecture that the thermal deconfining transition is governed by (e&eb) a universal mechanism, common to (e) all simple gauge groups.**

**Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 03 (2015) 057 DOI: 10.1007/JHEP03 (2015)057 Report number: DESY-14-146; IFT-UAM/CSIC-14-076 Cite as: arXiv: 1409.8305 [hep-lat] (or arXiv: 1409.8305v2 [hep-lat] for this version Exceptional thermodynamics: The equation of state of G(2) gauge theory Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini STOP In U (1) lattice gauge theory in three spacetime dimensions, the problem of confinement can be studied analytically in (e&eb) a semi-classical approach, in terms of (e&eb) a gas of monopoles with (e&eb) Coulomb-like interactions. In addition, this theory can be mapped to (e&eb) a spin model via (e&eb) an exact duality transformation, which allows (eb) one to perform high-precision numerical studies of (e) the confining potential. Taking advantage of these properties, authors carried out an accurate investigation of (e) the effective string describing (eb) the low-energy properties of flux tubes in (eb) this confining gauge theory. Authors found striking deviations from (e) the expected Nambu-Goto-like behavior, and, for the first time, evidence for (e) contributions that can be described by a term proportional to the extrinsic curvature of the effective string worldsheet. Such term is allowed by (e) Lorentz invariance and its presence in the infrared regime of the U (1) model was indeed predicted by (e) Polyakov several years ago. Results show that this term scales as expected according to (e) Polyakov's solution, and becomes (=) the dominant contribution to (e) the effective string action in the continuum limit. Authors also demonstrate analytically that the corrections to (e) the confining potential induced by (e) the extrinsic curvature term can be related to (e&eb) the partition function of (e) the massive perturbation of (e) a c=1 bosonic conformal field theory. The implications of results for SU (N) Yang-Mills theories in (e&eb) three and in four spacetime dimensions are discussed. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th) Journal reference: JHEP 1501 (2015) 105 DOI: 10.1007/JHEP01 (2015)105 Report number: IFT-UAM/CSIC-14-052 Cite as: arXiv:**

**1406.5127 [hep-lat] (or arXiv: 1406.5127v2 [hep-lat] for this version) A different kind of string Michele Caselle, Marco Panero, Roberto Pellegrini, Davide Vadacchino STOP In the past ten years, the ideas of supersymmetry have been profitably applied to (e&eb) many nonrelativistic quantum mechanical problems. In particular, there is now a much deeper understanding of why certain potentials are analytically solvable and (e&eb) an array of powerful new approximation methods for (e) handling potentials which are not (e) exactly solvable. In this report, authors review the theoretical formulation of (e) supersymmetric quantum mechanics and discuss many applications. Exactly solvable potentials can be understood in terms of (e&eb) a few basic ideas which include (e) supersymmetric partner potentials (e&eb) shape invariance and (e&eb) operator transformations. Familiar solvable potentials all have (e) the property of shape invariance. Authors describe new exactly solvable shape invariant potentials which include (e) the recently discovered self-similar potentials as (=) a special case. The connection between inverse scattering,(e&eb) isospectral potentials and (e&eb) supersymmetric quantum mechanics is discussed and multisoliton solutions of (e) the K dV equation are constructed. Approximation methods are also discussed within (eb) the framework of supersymmetric quantum mechanics and in particular it is shown that a supersymmetry inspired WKB approximation is (=) exact for a class of shape invariant potentials. Supersymmetry ideas give (eb) particularly nice results for the tunneling rate in (eb) a double well potential and for improving (eb+) large N expansions. Authors also discuss the problem of (e) a charged Dirac particle in (eb) an external magnetic field and other potentials in terms of (e&eb) supersymmetric quantum mechanics. Finally, authors discuss structures more general than supersymmetric quantum mechanics such as(=) parasupersymmetric quantum mechanics in which there is symmetry between a boson and (e&eb) a para-fermion of order p. Copyright © 1995 Published by Elsevier B.V. Physics Reports Volume 251, Issues 5–6, January 1995, Pages 267–385Supersymmetry and quantum mechanics Fred Coopera, Avinash Khareb, Uday Sukhatmec STOP The free energy of (e) the maximally**

**supersymmetric SU (N) gauge theory at temperature T is expected to scale, in the large-N limit, as (=) N2T4 times a function of the 't Hooft coupling, f (gYM2N). In the strong coupling limit the free energy has been deduced from (e) the near-extremal 3-brane geometry, and its normalization has turned out to be Full-size image (<1 K) times that found in (eb) the weak coupling limit. In this paper authors calculate the leading correction to (e) this result in inverse powers of (e) the coupling, which originates from (e) the R4 terms in the tree-level effective action of (e) type IIB string theory. The correction to Full-size image (<1 K) is positive and of (e) order Full-size image (<1 K). Thus, f (gYM2N) increases as (e&eb) the 't Hooft coupling is decreased, in accordance with (e&eb) the expectation that is should be approaching 1 in (eb) the weak coupling limit. Authors also discuss similar corrections for (e) other conformal theories describing (eb) coincident branes. In particular, authors suggest (eb) that the coupling independence of (e) the near extremal entropy for (e) D1-branes bound to (e) D5-branes is related to (e&eb) the vanishing of (e) the Weyl tensor of AdS3 × S3. Copyright © 1998 Published by Elsevier B.V. Nuclear Physics B Volume 534, Issues 1–2, 23 November 1998, Pages 202–222Coupling constant dependence in the thermodynamics of N = 4 supersymmetric Yang-Mills theory Steven S. Gubsera, Igor R. Klebanovb, 1, Arkady A. Tseytlinc, 2 STOP A new approach to (e) N=2 supersymmetry based on (e) the concept of harmonic superspace is proposed and is used to give (eb) an unconstrained superfield geometric description of N=2 super Yang-Mills and supergravity theories as (=) well as of matter N=2 hypermultiplets. The harmonic N=2 superspace has (e) an independent coordinates, in addition to the usual ones, the isospinor harmonics ui+or- on (e&eb) the sphere SU (2)/U (1). The role of ui+or- is to relate the SU (2) group realised on the component fields to (e&eb) a U (1) group acting on (e&eb) the relevant superfields. Their introduction makes it possible to SU (2)-covariantise (eb (t), e&eb) the notion of Grassmann analyticity. Crucial for the construction is the existence of (e) an analytic subspace of the general harmonic N=2 superspace. The hypermultiplet superfields and (e&eb) the true prepotentials (pre-**

**prepotentials) of N=2 super Yang-Mills and supergravity are (=) unconstrained superfunctions over (e&eb) this analytic subspace. The pre-prepotentials have (e) a clear geometric interpretation as gauge connections with (e&eb) respect to the internal SU (2)/U (1) directions. A radically new feature arises: the number of gauge and (e&eb) auxiliary degrees of freedom becomes (=) infinite while (e) the number of physical degrees of freedom remains (=) finite. Other new results are the massive N=2 Yang-Mills theory and (e&eb) various off-shell self-interactions of hypermultiplets. The propagators for matter and Yang-Mills superfields are given. Unconstrained N=2 matter, Yang-Mills and supergravity theories in harmonic superspace A Galperin, E Ivanov, S Kalitzin, V Ogievetsky and E Sokatchev STOP By assuming the existence of a novel multipronged string state for (e) D-particles interacting with (e&eb) D-brane intersections in type IIA string theory, authors are able to derive a quantum mechanical description of (e) supersymmetric Reissner-Nordstrom black holes. A supersymmetric index calculation provides (eb) evidence for (e) this conjecture. The quantum mechanical system becomes (=) two decoupled conformal quantum mechanical systems in (eb) the low energy limit. The conformal quantum mechanics has expected properties of (e) a dual description of string theory on (e&eb) AdS2×S2. Subjects: High Energy Physics - Theory (hep-th) Journal reference: JHEP 0001:006, 2000 DOI: 10.1088/1126-6708/2000/01/006 Report number: LBNL-43734 Cite as: arXiv:hep-th/9907101 (or arXiv:hep-th/9907101v5 for this version) Supersymmetric Quantum Mechanical Description of Four Dimensional Black Holes Julie D. BlumSTOP Using recently proposed soliton equations authors derive a basic identity for the scaling violation of (e) N=2 supersymmetric gauge theories Σiai∂F/∂ai−2F=8πib1u. Here F is (=) the prepotential, ai's are (=) the expectation values of (e) the scalar fields in (eb) the vector multiplet, u=1/2 Tr <ϕ2> and b1 is the coefficient of (e) the one-loop β-function. This equation holds in (eb) the Coulomb branch of all N=2 supersymmetric gauge theories coupled with (e&eb) massless matter. TOHRU EGUCHI and SUNG-KIL YANG, Mod Phys. Lett. A 11, 131 (1996).**

**DOI: 10.1142/S0217732396000151 PREPOTENTIALS OF N=2 SUPERSYMMETRIC GAUGE THEORIES AND SOLITON EQUATIONS STOP The Coulomb branch of N = 2 supersymmetric gauge theories in four dimensions is described in general by (e) an integrable Hamiltonian system in (eb) the holomorphic sense. A natural construction of such systems comes from (e) two-dimensional gauge theory and (e&eb) spectral curves. Starting from this point of view, authors propose an integrable system relevant to (e&eb) the N = 2 SU (n) gauge theory with (e&eb) a hypermultiplet in the adjoint representation, and offer (eb) much evidence that it is correct. The model has (e) a Full-size image (<1 K) S-duality group (with the central element −1 of Full-size image (<1 K) acting as (=) charge conjugation); Full-size image (<1 K) permutes the Higgs, confining, and oblique confining phases in the expected fashion. We also study more exotic phases. Copyright © 1996 Published by Elsevier B.V. Nuclear Physics B Volume 460, Issue 2, 5 February 1996, Pages 299–334Supersymmetric Yang-Mills theory and integrable systems Ron Donagia, 1, Edward Wittenb, 2 STOP A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 present a first step towards generalizing (eb) the work of Seiberg and Witten on N = 2 supersymmetric Yang-Mills theory to (e&eb) arbitrary gauge groups. Specifically, authors propose a particular sequence of (e) hyperelliptic genus n − 1 Riemann surfaces to underly (e&eb) the quantum moduli space of SU (n)N = 2 supersymmetric gauge theory. These curves have (e) an obvious generalization to (e) arbitrary simply laced gauge groups, which involves (e&eb) the A-D-E type simple singularities. To support proposal, authors argue that the monodromy in (e&eb) the semiclassical regime is (=) correctly reproduced. We also give some remarks on a possible relation to string theory. Copyright © 1995 Published by Elsevier B.V. Physics Letters B Volume 344, Issues 1–4, 26 January 1995, Pages 169–175Simple singularities and N = 2 supersymmetric Yang-Mills theory A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 STOP Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 review and extend earlier work that uses (e) the AdS/CFT correspondence to relate the**

**black-hole–black-string transition of gravitational theories on a circle to (e&eb) a phase transition in maximally supersymmetric (1 + 1)-dimensional SU(N) gauge theories at (eb) large N, again compactified on (e&eb) a circle. Authors perform gravity calculations to determine (eb) a likely phase diagram for (e) the strongly coupled gauge theory. Authors then directly study the phase structure of (e) the same gauge theory, now at (eb) weak't Hooft coupling. In the interesting temperature regime for (e) the phase transition, the (1 + 1)-dimensional theory reduces to (e&eb) a (0 + 1)-dimensional bosonic theory, which they solve using (e) Monte Carlo methods. Authors find strong evidence that the weakly coupled gauge theory also exhibits (eb) a black hole–black string-like phase transition in (eb) the large N limit. Authors demonstrate that a simple Landau–Ginzburg-like model describes (eb) the behaviour near the phase transition remarkably well. The weak coupling transition appears to be (=) close to the cusp between a first-order and a second-order transition. Black-hole–black-string phase transitions in thermal (1 + 1)-dimensional supersymmetric Yang–Mills theory on a circle Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 Published 29 October 2004 • 2004 IOP Publishing Ltd Classical and Quantum Gravity, Volume 21, Number 22STOP Monitoring the orbits of stars around (e&eb) Sgr A* offers(eb)the possibility of detecting (e&eb) the precession of (e) their orbital planes due to (e) frame dragging, of (e) measuring the spin and quadrupole moment of (e) the black hole, and of (e)testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and (e&eb) tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are (=), in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for which an optical interferometer, such as GRAVITY, will detect orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to (e&eb) the timescale of precession due to (e) the quadrupole moment of (e) the black hole. As a result, this non-conservative effect is (=) a**

**potential source of systematic uncertainty in (eb) testing the no-hair theorem with (e&eb) stellar orbits. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI:10.1088/0004-637X/777/1/57 Cite as:arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this version) Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP Authors perform a phase space analysis of (e) a generalized modified gravity theory with (e&eb) non-minimal coupling between geometry and (e&eb) matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that (eb) in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, Authors prove that there is (=) an attractor solution to describe (eb) the late time accelerating universe when (e) the modified gravity is chosen in (eb) a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and supernova type Ia data sets, Authors find that (eb) the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as:arXiv: 1512.09281 [gr-qc] (or arXiv: 1512.09281v1 [gr-qc] for this version) Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang study the P−V criticality and phase transition in (eb) the extended phase space of (e) charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of the system and (e&eb) its conjugate quantity is (=) the thermodynamic volume of (e) the black hole. The black**

**holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb)no P−V criticality and phase transition appear, while for the black holes with a spherical horizon, even when (e) the charge of the black hole is absent, the P−V criticality and the small black hole/large black hole phase transition (eb) will appear, but it happens only in (eb) d=5 dimensions; when (e) the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to (e) have the P−V criticality and the small black hole/large black hole phase transition, there exists (eb) an upper bound for (e) the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the Gauss-Bonnet coefficient and Q is the charge of the black hole. Authors calculate the critical exponents at the critical point and find that for all cases, they are (=) the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09 (2013)005 Cite as: arXiv: 1306.6233 [gr-qc] (or arXiv: 1306.6233v4 [gr-qc] for this version) P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP paper studies the whole process of (e) gravitational collapse and (e&eb) accretion of (e) a massless scalar field in (eb) asymptotically flat spacetime. Two kinds of initial configurations are considered. One is the initial data without (e) black hole; the other contains (e) a black hole. Under suitable initial conditions, Authors find that multi-horizon will appear which mean (eb) that the initial black hole formed by (e) gravitational collapse or existing at the beginning will instantly expand and suddenly grow (e&eb, eb), rather than grows gradually in the accretion process. A new type of critical behavior is found around (e&eb) the instant expansion. The numerical computation shows (eb) that the critical exponents are (=) universal. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Computational Physics (physics. Comp-ph) Cite as: arXiv: 1604.03363 [gr-qc] (or arXiv: 1604.03363v1 [gr-qc] for this version)**

**Multi-horizon and Critical Behavior in Gravitational Collapse of Massless Scalar Zhoujian Cao, Rong-Gen CAI, Run-Qiu Yang STOP In this paper, Authors present the results on the gravitational collapse of (e) charged massless scalar field in asymptotically flat spacetime with(e&eb) a perfectly reflecting wall. Differing from previous works, Authors study the system in (eb) the double null coordinates, by which we could simulate (e&eb) the system until the black hole forms with (e&eb) higher precision but less performance time. Authors investigate the influence of charge on (e&eb) the black hole formation and the scaling behavior near the critical collapses. The gapless and gapped critical behaviors for (e) black hole mass and charge are studied numerically. Authors find that they satisfy (eb) the scaling laws for critical gravitational collapse but the gapped critical behavior is different from (e) its AdS counterpart. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1602.00112 [gr-qc] (or arXiv: 1602.00112v1 [gr-qc] for this version) Multiple critical gravitational collapse of charged scalar with reflecting wall Rong-Gen CAI, Run-Qiu Yang STOP In this paper Authors investigate the (2+1)-dimensional topological non-Fermi liquid in (eb) strongly correlated electron system, which has (e) a holographic dual description by (e) Einstein gravity in (3+1)-dimensional anti-de Sitter (AdS) space-time. In a dyonic Reissner-Nordstrom black hole background, Authors consider a Dirac fermion coupled to (e&eb) the background U (1) gauge theory and an (e&eb) intrinsic chiral gauge field bM induced by (e) chiral anomaly. UV retarded Green's function of (e) the charged fermion in (eb) the UV boundary from (e) AdS4 gravity is calculated, by imposing (e&eb) in-falling wave condition at the horizon. Authors also obtain IR correlation function of (e) the charged fermion at (eb) the IR boundary arising from (e) the near horizon geometry of (e) the topological black hole with (e&eb) index k=0,±1. By using (e) the UV retarded Green's function and IR correlation function Authors analyze the low frequency behavior of (e) the topological non-Fermi liquid at zero and (e&eb) finite temperatures, especially the relevant non-Fermi liquid behavior near (eb) the quantum critical region. In addition, Authors**

**find that when k=±1, the effective mass of (e) bulk Dirac fermion at IR fixed point is (=) topologically quantized, consequently the Fermi momentum presents (eb) a discrete sparse distributed pattern, in contrary to (e&eb) the Ricci flat case with k=0. As demonstrated examples, Authors calculate the spectral functions of (e) topological non-Fermi liquid due to (e) anomalous Hall effect (AHE) (|b→|≠0,b 0=0) and chiral magnetic effect (CME) (|b→|=0,b0≠0), which can be tested in topological materials such as (=) topological insulator and (e&eb) topological metal with (e&eb) strong couplings. Subjects: High Energy Physics - Theory (hep-th); Strongly Correlated Electrons (cond-mat.str-el); General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv: 1601.03865 [hep-th] (or arXiv: 1601.03865v1 [hep-th] for this version) Topological Non-Fermi Liquid Rong-Gen Cai, Yong-Hui Qi, Yue-Liang Wu, Yun-Long Zhang STOP The unusual magnetic materials are significant in both science and technology. However, because of (e) the strongly correlated effects, it is difficult to understand their novel properties from (e) theoretical aspects. Holographic duality offers (eb) a new approach to understanding such systems from (e) gravity side. This paper will give a brief review of Authors recent works on the applications of holographic duality in understanding (e&eb) unusual magnetic materials. Some quantitative compare between holographic results and (e&eb) experimental data will be shown and some predictions from (e) holographic duality models will be discussed. Comments: For Second LeCosPA Symposium at Taiwan National University, Taipai, Dec. 14-18, 2015 Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1601.02936 [hep-th] (or arXiv:1601.02936v1 [hep-th] for this version) Understanding strongly coupling magnetism from holographic duality Rong-Gen Cai, Run-Qiu Yang STOP It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a**

**Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3**

**[astro-ph.CO] for this version) Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang**

# Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

## Module Number Six

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-465) 1990. We demonstrate its applicability cosmology/ quantum computation economics etal. For the first time, physicists couple distant nuclear spins using a single electron July 11, 2016. Some sentences have been deleted due to spatial restrictions. Please excuse me. For the first time, researchers at the University of Basel in Switzerland have coupled the nuclear spins of distant atoms using (e) just a single electron. Three research groups took part in this complex experiment, the results of which have now been published in the journal Nature Nanotechnology. In most materials, the nuclear spins of (e) neighboring atoms have (e) only a very weak effect on one another, as (=) the tiny nuclei are located deep within (e&eb) the atoms. This is not the case in metals, however, where (e) some of the electrons are able to move freely. The electron spins are able to couple nuclear spins at (eb) relatively large distances from one another. Named after four physicists, this RKKY interaction was discovered in the 1950s. Nuclear spins linked by (e&eb) an individual electron spin Now, an experiment by researchers from the Department of Physics at the University of Basel has for the first time succeeded in demonstrating this mechanism with (e&eb) a single electron, describing it using (e) quantum theory. The team led by Prof. Richard Warburton trapped a single electron inside (eb) a quantum dot. With use of a method developed in Basel to measure (eb) nuclear spin resonance, they showed that the electron coupled (e&eb) the spins of nuclei up to five nanometers apart - a huge distance in the world of (e) nuclear spins. The results are particularly relevant to (e&eb) the development of spin qubits: these seek to use (e) electron spins to carry (e&eb) information, but the interaction with (e&eb) the nuclei limits (e) the stability of the quantum information. Concentrated blast of physics More information: Role of the electron spin in determining the coherence of the nuclear spins in a quantum dot, Nature Nanotechnology, DOI: 10.1038/nnano.2016.114 Journal reference: Nature Nanotechnology search and more info website Provided by: University of Basel Read more at: http://phys.org/news/2016-07-physicists-couple-distant-nuclear-electron.html#jCp A new spin on quantum computing: Scientists train electrons with microwaves February 15, 2016 A silicon sample doped with (e&eb) bismuth atoms (left image) that is (=) just 150 nanometers thick is fitted with (e&eb) a superconducting resonator that includes (e) a capacitor (black, in left image; light gray in center image) and an inductive wire (red line … In what may provide (eb) a potential path to processing information in (eb) a quantum computer, researchers have switched an intrinsic property of electrons from an excited state to (e&eb) a relaxed state on demand using (e) a device that served as (=) a microwave "tuning fork." The team's findings could also lead to (eb) enhancements in magnetic resonance techniques, which are (=) widely used to explore (e) the structure of materials and biomolecules, and for medical imaging. The**

**international research team, which included scientists at the Department of Energy's Lawrence Berkeley National Laboratory (Berkeley Lab), demonstrated how to dramatically increase the coupling of (e) microwaves in a specially designed (e&eb) superconducting cavity to (e) a fundamental electron property called spin—which, like a coin, can be flipped. By zapping an exotic silicon material developed at Berkeley Lab with the microwaves, they found that they could rapidly change (e&eb) the electron spins from (e) an excited state to a relaxed, ground state by (e) causing (eb) the electrons to emit (eb) some of their energy in the form of (e) microwave particles known as photons. Left on their own, the electron spins would be extremely unlikely to flip back to (e) a relaxed state and to also emit (eb) a photon - the natural rate for this light-emitting effect, known as the Purcell effect, is about once every 10,000 years. The experiment demonstrated an accelerated, controllable relaxation of (e) electron spins and the release of (e) a microwave photon in about 1 second, said Thomas Schenkel, a physicist in Berkeley Lab's Accelerator Technology and Applied Physics Division who led the design and development of the silicon-bismuth sample used in the experiment. A new spin on quantum computing: Scientists train electrons with microwaves Like the classical tuning fork slung from guitar strings in this photo, scientists used (e) a device that worked like a microwave tuning fork to switch (e&eb) a fundamental property of electrons on (e&eb) demand in a silicon-bismuth sample developed at Berkeley Lab. Credit: Flickr/Pierre Guinoiseau "It's like a juggler who throws the balls up, and the balls come down 1,000 times faster than normal, and they also emit (eb) a microwave flash as they drop," he said. The results were published online Feb. 15 in the journal Nature. "Our results are highly significant for (e&eb) quantum information processing," said Patrice Bertet, a quantum electronics scientist at the French Atomic Energy Commission (CEA) who led the experiment. "Indeed, they are a first step toward (e&eb) the strong coupling of (e) individual electron spins to microwave photons, which could form (eb) the basis of new spin-based quantum computer architecture." John Morton, a professor at the London Center for Nanotechnology and co-author of the study, said, "Our ultimate aim is to find a link between quantum information that is fixed and (e&eb) quantum information that can be transported by photons." In today's computers, information is stored as (=) individual bits, and each bit can either be (=) a one or a zero. Quantum computers, though, could conceivably be (=) exponentially more powerful than modern computers because (e) they would use a different kind of bit, called a qubit, which because of (e) the weird ways of quantum mechanics can simultaneously behave (e&eb) as both a one and a zero. A coupled array of qubits would allow (eb) a quantum computer to perform (e&eb) many, many calculations at (eb) the same time, and electron spins are (=) candidates for qubits in (eb) a quantum computer. The latest study shows how the microwave photons could work in concert with (e&eb) the spins of electrons to move information in a new type of computer. "What we need now is ways to wire up these systems—to couple (e&eb) these spins together," Morton said. "We need to make coupled qubits that can perform (e&eb) computations." In the experiment, conducted at CEA in France, a small sample of a highly purified form of (e) silicon was implanted with (e&eb) a matrix of bismuth atoms, and a superconducting aluminum circuit was deposited on (e&eb) top to create (eb) a high-quality resonant cavity that allowed (eb) precise tuning of the microwaves. The electron spins of the bismuth atoms were then flipped into (e&eb) the excited, "spin-up" state. The microwave cavity was then tuned (e&eb) like a musical tuning fork, to (e) a particular resonance that coaxed the spins into emitting (eb) a photon as they flipped back to (e) a relaxed state. The cavity boosted the number of states into (e&eb) which a photon can be (=) emitted, which greatly increased (eb) the decay rate for (e) the electron spins in (eb) a controllable way. The technique is much like buying more lottery tickets to increase (eb) your chances of winning, Morton said. A new spin on quantum computing: Scientists train electrons with microwaves From left are, Berkeley Lab scientists Thomas Schenkel, Qing Ji and Peter Seidl at the NDCX-II (Neutralized Drift Compression Experiment II), which produces (eb) powerful ion beams. Researchers are exploring how to use NDCX-II to process exotic … The large bismuth atoms embedded in (eb) the silicon sample provided (eb) the electrons with (e&eb) unique spin properties that enabled (eb) the experiment. Schenkel said that implanting the bismuth atoms into (e&eb) the delicate silicon framework, a**

**process known as "doping," was "like squeezing bowling balls into (e&eb) a lattice of ping-pong balls." "We did a new trick with silicon. People wouldn't expect you could squeeze (e) anything new out of (e) silicon," Schenkel said. "Now we're looking into further improving (eb) bismuth-doped silicon and into (e&eb) tailoring the spin properties of (e) other materials, and using (e) this experimental technique for these materials." To enhance the performance of materials used (e) in future experiments, Schenkel said it will be necessary to improve (eb) the doping process so it is less damaging to (e) the silicon lattice. Also, the implantation process could be designed to produce (eb) regularly spaced arrays of (e) individual electron spins that would be more useful for (e) quantum computing than a concentrated ensemble of (e) electron spins. "We are now doing experiments on processing (e&eb) this and other materials at higher temperature and pressure with nanosecond ion pulses at NDCX-II, one of the accelerators here at Berkeley Lab," Schenkel said. "There are indications that it will improve (eb) the overall spin quality." Researchers said the latest research could potentially prove useful in boosting (eb) the sensitivity of scientific techniques like nuclear magnetic resonance spectroscopy and dynamic nuclear polarization, useful for a range of experiments, and could also shorten (e) experimental times by manipulating (e&eb) spin properties. "You need a way to reset (e&eb) spins—the ability to cause (eb) them to relax on demand to improve (eb) the rate at which you can repeat an experiment," Morton said. Bertet said it may be possible to further accelerate (eb) the electron-flipping behavior to (e) below 1 millisecond, compared to (e) the 1-second rate in the latest results. "This will then open the way to many new applications," he said. Explore further: Bismuth provides perfect dance partners for quantum computing qubits More information: A. Bienfait et al. Controlling spin relaxation with a cavity, Nature (2016). DOI: 10.1038/nature16944 Nature | Letter Controlling spin relaxation with a cavity A. Bienfait, J. J. Pla, Y. Kubo,X. Zhou, M. Stern, C. C. Lo, C. D. Weis, T. Schenkel, D. Vion,D. Esteve, J. J. L. Morton& P. Berte 531, 74–77 (03 March 2016) doi:10.1038/nature16944 Spontaneous emission of radiation is (=) one of the fundamental mechanisms by which an excited quantum system returns to (e) equilibrium. For spins, however, spontaneous emission is (=) generally negligible compared to (e&eb) other non-radiative relaxation processes because of (e) the weak coupling between the magnetic dipole and (e&eb) the electromagnetic field. In 1946, Purcell realized that the rate of spontaneous emission can be greatly enhanced by (e) placing the quantum system in a resonant cavity. This effect has since been used extensively to control (e&eb) the lifetime of atoms and semiconducting heterostructures coupled to (e&eb) microwave2 or optical3, 4 cavities, and is essential for (e) the realization of high-efficiency single-photon sources5. Here authors report the application of this idea to (e&eb) spins in solids. By coupling donor spins in silicon to (e&eb) a superconducting microwave cavity with (e&eb) a high quality factor and a small mode volume, authors reach the regime in which spontaneous emission constitutes (e) the dominant mechanism of spin relaxation. The relaxation rate is increased by (e) three orders of magnitude as the spins are tuned to (e) the cavity resonance, demonstrating that energy relaxation can be controlled (e&eb) on demand. Results provide a general way to initialize (eb) spin systems into their ground state and therefore have applications in (eb) magnetic resonance and (e&eb) quantum information processing6. They also demonstrate that the coupling between the magnetic dipole of a spin and (e&eb) the electromagnetic field can be enhanced up to (e) the point at which quantum fluctuations have (e) a marked effect on (e&eb) the spin dynamics; as such, they represent (eb) an important step towards the coherent magnetic coupling of (e) individual spins to (e) microwave photons Scientists obtain evidence of many-body localization in a closed quantum system July 8, 2016 by Olivia Meyer-Streng Science Cover Vol 352, Issue 6293, 24 June 2016, to highlight the research published in this issue (see original publication). The cover shows an artist's view of the experiment, illustrating the persistence of the density step of (e) an atomic distribution in an optical lattice. During equilibration ordinary many-body systems lose (e) all information about (e&eb) the initial state. Every morning authors experience an example for (e) this behaviour. Milk poured into a cup of coffee mixes (e&eb) perfectly and after some time it is impossible to say how exactly the two fluids were (=) put together. The same behaviour holds for (e) almost all quantum systems. However,**

**recently a new phenomenon called "many-body localization" has been predicted theoretically, which allows (eb) well insulated quantum systems to (e) preserve memory of the initial state forever. Now a team of scientists around Dr. Christian Groß and Professor Immanuel Bloch (Director at MPQ and Chair of Quantum Optics at LMU Munich), in cooperation with David Huse (Princeton University), has obtained evidence of such a behaviour in a two-dimensional quantum system of cold rubidium atoms trapped in (e&eb) an optical lattice. The scientists observed that – beyond a certain degree of disorder imprinted on (eb) the particle ensemble in the beginning – the system would relax into (e&eb) a steady state still containing (e) detailed microscopic information about its past. "We were able to observe the transition from a thermalized state into (e&eb) a many-body localized phase", Christian Groß points out. "It is the first observation of that kind in a regime that is not accessible with (e&eb) state-of-the-art simulations on classical computers." The experiment is not only of fundamental interest; the results might also lead to (eb) new ways for storing quantum information. Motivated by the foundational problem of how interacting particles behave (e&eb) in a disordered system, in the 1950s the American physicist Philip Warren Anderson discovered the famous localization phenomenon for (e) non-interacting particles, now called "Anderson localization". Here, disorder prevents (e) the particles to move and consequently all transport is (=) stopped. But what happens when disorder comes together with (e&eb) interactions? Will interactions lead to (eb) transport and thermalization, or will the localization persist even at (eb) high energies? So far, there is no theoretical model that faithfully predicts (eb) the evolution of a closed quantum system in more than one dimension under (e&eb) these conditions, although, the possibility for (e)localization has been theoretically suggested. In order to investigate these questions experimentally, strict requirements on controllability and perfect isolation of (e) the system have to be fulfilled. In the experiment described here, ultracold rubidium atoms are loaded into (eb) an optical lattice, a microscopic array of light traps formed by interfering (e&eb) laser beams. The disorder is created by projecting (e&eb) a computer-generated random light pattern onto (e&eb) the optical lattice. As a result of this "disorder", the depth of each microscopic trap in the lattice potential varies from site to (e) site. In fact, Prof. Bloch's group has advanced the experimental tools to such perfection that they can steer the position of (e) the atoms in the artificial light crystal and (e&eb) the interaction between (e&eb) them almost at will. With a high resolution microscope that detects the emitted fluorescence light, the position of each atom can be observed with (e&eb) high precision. Additionally, the atomic density distribution of the initial state can be (=) controlled, and for variable time intervals the evolution of the distribution can be (=) measured with high precision. These tools at hand, non-thermalizing behaviour can be probed in (eb) a conception ally very simple way. Any thermalized state of a closed system reflects (e&eb) the symmetry of its container – like, for example, water spilled into (e&eb) a round bowl would immediately cover (e) the whole bottom. In analogy, the scientists generate (eb) a density step in the initially prepared sample by (e) blowing away half of the atomic distribution with laser radiation. Then they watch how the remaining particles migrate into (e&eb) the empty half. For small disorders, the initially prepared density step is smeared out (e) fast and the initially empty and filled halves become (=) indistinguishable. However, when the measurement is repeated for (e) strong disorder, traces of the initial state remain and the system does not (e) relax to a thermal state even for (e) very long times. "We observe a fairly sharp onset of non-thermalizing behaviour above (e) a critical value", says Christian Groß. "This absence of thermalization is remarkable because (e) thermalization persists in a system of interacting particles, even at (eb) the high energies probed in the experiment." The scientists interpret their observation as (=) the onset of many-body localization in (eb) the atomic system. This is of fundamental interest because it means (eb) the breakdown of equilibrium statistical mechanics. On the other hand, the persistence of initial state information could be used as (=) a source for quantum information technologies. "It should also be emphasized that we obtain these results for a system size that is far beyond (e) numerically accessible scales", says Jae-yoon Choi, postdoc at the experiment. More information: J.-y. Choi et al. exploring the many-body localization transition in two dimensions, Science (2016). DOI: 10.1126/science.aaf8834 Journal reference: Science search and more**

**info website Provided by: Max Planck Society search and more info website Bosons refusing to thermalize in 2D Messy, interacting quantum-mechanical systems are difficult to analyze theoretically. In a single spatial dimension, the calculations are (=) still tractable, and experiments have recently confirmed (eb) the prediction that sufficiently strong disorder can disrupt (e&eb) the transport of interacting particles. In two dimensions, however, the theoretical blueprint is missing. Choi et al. used (e) single-site imaging of cold 87Rb atoms in (eb) an optical lattice to show (eb) that similar localization occurs in (eb) two-dimensional (2D) systems. The study highlights the power of quantum simulation to solve problems that are currently inaccessible to classical computing techniques. Science, this issue p. 1547 Abstract A fundamental assumption in statistical physics is that generic closed quantum many-body systems thermalize under (e&eb) their own dynamics. Recently, the emergence of many-body localized systems has questioned (e) this concept and challenged our understanding of the connection between statistical physics and (e&eb) quantum mechanics. Here authors report on the observation of a many-body localization transition between thermal and (e&eb) localized phases for (e) bosons in a two-dimensional disordered optical lattice. With our single-site–resolved measurements, authors track the relaxation dynamics of (e) an initially prepared out-of-equilibrium density pattern and find strong evidence for (e) a diverging length scale when (e) approaching the localization transition. Experiments represent a demonstration and in-depth characterization of many-body localization in (eb) a regime not accessible with (e&eb) state-of-the-art simulations on classical computers. Pelican instrument provides crucial experimental evidence of unusual quantum state July 8, 2016 Artist's illustration. Credit: Australian Nuclear Science and Technology Organisation (ANSTO) Experiments using inelastic neutron scattering at the Australian Centre for Neutron Scattering have found indications (eb) of a possible new quantum spin state in a novel antiferromagnetic material barium ytterbium zinc oxide (Ba3Yb2Zn5O11) which provides (eb) both a challenge and validation of (e) the third law of thermodynamics. The quantum mechanical information gained from (e) the experiment is crucial not only in understanding the quantum state of the specific material but also provides (eb) better understanding of high temperature superconductivity and ultimately offers (eb) guidance in the design of materials for (e) quantum computing. A research team led by Associate Professor Takatsugu Masuda at the Institute for Solid State Physics at the University of Tokyo with collaborators from ANSTO, Osaka University and Taiwan's National Cheng Kung University, have performed a successful experiment using Pelican, time-of-flight cold neutron spectrometer at ANSTO. The results have just been published in Physical Review B as a rapid Communication and an Editors' Suggestion. "The great success of this challenging experiment relied on (e) the excellent performance of the Pelican instrument, wonderful technical support at ANSTO and strong expertise in the relevant field from A/Prof Masuda's group," said Dr Dehong Yu and Dr Richard Mole, co-authors and Pelican instrument scientists. Pelican instrument provides crucial experimental evidence of unusual quantum state INS spectra measured at 1.5 K (a), 6 K (b), 12 K (c), and 40 K (d) using setup I. (e) INS spectrum measured at 1.5 K using the setup II. (f) Q dependence of the integrated intensity obtained from spectrum in (e). Red solid curve is the …….(deleted due to spatial constraints) Barium ytterbium zinc oxide (Ba3Yb2Zn5O11) is a novel antiferromagnetic model system having (e) a "breathing pyrochlore lattice" consisting of (e) arrays of alternating large and (e&eb) small corner-sharing tetrahedra of ytterbium atoms. Earlier research had suggested (eb) the material was (=) an isotropic spin half (S= ½) system. However, four low energy excitations measured on the sample with the Pelican instrument at a temperature of 1.5 K clearly disagree with (e) this designation and the associated theoretical model based on (e) the isotropic spins having simple interactions. "These observations suggested (eb) a much more complex spin state in the system," said Yu. "The experiment is interesting and important because as you start going very low in temperature, quantum mechanics takes over (e) system rather than thermodynamics, and you observe unusual properties," said Mole. "According to low temperature heat capacity measurements, the material does not (e) have zero entropy or a single ground state until you reach 63 milli Kelvin or -272.94 Celsius, before it starts obeying the third law." Pelican instrument provides crucial experimental evidence of (e) unusual quantum state**

**Breathing pyrochlore lattice formed by (e) A3+ ions as determined by Kimura, Nakatsuju and Kimura in previous work The third law of thermodynamics states that for pure crystals as the temperature of a system approaches absolute zero K, (eb) its entropy or disorder approaches zero. "The research clearly indicates (eb) that below this temperature the material may be in (eb) an exotic state, such as (=) a liquid-like spin state, in which the arrangement of spins behaves (e&eb) like a liquid. In this exotic state, novel phenomena may show up and be detectable; for example, charge and spin separation of electrons may be (=) possible," said Yu. The researchers suggest that further investigations, experimentally and theoretically, are required to ultimately determine the true identity of (e) the ground state of the material. "Discovering the existence of regular tetrahedral spin made it well worth coming all the way to Australia to make use of this neutron experimental apparatus," said Associate Professor Masuda. Many Japanese researchers are using ANSTO's neutron scattering instruments through the peer review system, while their nuclear facilities are offline. The outcome of this particular experiment is also part of the collaboration agreement between ANSTO and the University of Tokyo. More information: T. Haku et al. Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet, Physical Review B (2016). DOI: 10.1103/PhysRevB.93.220407 Journal reference: Physical Review B search and more info website: provided by: Australian Nuclear Science and Technology Organisation (ANSTO) search and more info website Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 T. Haku, K. Kimura, Y. Matsumoto, M. Soda, M. Sera, D. Yu, R. A. Mole, T. Takeuchi, S. Nakatsuji, Y. Kono, T. Sakakibara, L.-J. Chang, and T. Masuda Phys. Rev. B 93, 220407(R) – Published 24 June 2016 Authors study low-energy excitations in (eb) the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 by (e) a combination of inelastic neutron scattering (INS) and (e&eb) thermodynamical property measurements. The INS spectra are quantitatively explained by (e) spin-12 single-tetrahedron model having (e) XXZ anisotropy and Dzyaloshinskii-Moriya interactions. This model has a twofold degeneracy of (e) the lowest-energy state per tetrahedron and well reproduces (eb) the magnetization curve at (eb) 0.5 K and heat capacity above 1.5 K. At lower temperatures, however, authors observe a broad maximum in (eb) the heat capacity around 63 mK, demonstrating (eb) that a unique quantum ground state is selected due to (e) extra perturbations with (e&eb) an energy scale smaller than the instrumental resolution of INS. Revised 13 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.220407 © 2016 American Physical Society Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA (Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor. In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface. As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone. Such planets are all the "right" distance from (e) their star, and hopefully, also have water. But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone. For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement. But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway. But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain. It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of**

**reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life. Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface. The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113 Abstract The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question. Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry. This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry. Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments. Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy. The electronic and (e&eb) structural variability among them is extraordinary. The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations. The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface. The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers. Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp**

## 12. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_6}{dt} = -b_6 Q_6 + g_6(g_6)Q_6$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_6(0) = 0, \quad g_6(P_6) > 0, \quad P_6 \ 0, \qquad \lim_{P_6 \to \infty} g_6(P_6) = (K_6)_6 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_6 + (\rho_a)_6 + D_6(a_6, P_6, Q_6)\rho_6 + 0, \qquad 0 < a_6 \ \infty,$$

$$\rho_6(t, 0) = \int_0^\infty n_6(a_6, P_6, Q_6)\rho_6(t_6, a_6)da_6$$

Where $D_6$ is the mortality function and $n_6$ is the fertility function.

We assume following equations held

$$\left[\frac{dQ_6}{dt}\right]_6 = x_6\left[\frac{dQ_2}{dt}\right]_6 - x_2\left[\frac{dQ_6}{dt}\right]_6$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_6 - x_4\left[\frac{dQ_2}{dt}\right]_6$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_6 - x_6\left[\frac{dQ_3}{dt}\right]_6$$

Where $x_6, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_6$ and $g_6(P_6)Q_6$

Now when you add the term $(b_6Q_6)$ with $g_6(P_6)(Q_6)$ we are infect see is that $(b_6Q_6)$ is credited incrementally part of $g_6(P_6)(Q_6)$ or instant nearly, when subtracted from $\frac{dQ_6}{dt}$ , it means that $g_6(P_6)(Q_6)$ is slowly eaten by $\frac{dQ_6}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_6 = \mu_6(a_6) + (f_6)_6(P_6)_6 + (f_2)_6(Q)_6 \quad (2)$$

$$n_6 = \beta_6(a_6)h_6(Q_6), \quad (3)$$

Where $\mu_6$the age- dependent debit rate whit is $\mu_6(a_6)$ 0 $\forall a > 0$; $(f_6)_6$ is the density- dependent death rate with $f_6(0) = 0$, $(f_6^6)_6(P_6)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_6)_6(P_6) = (k_6)_6$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_6(0) = 0, (f_2^6)_6(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_6(Q_6) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 6, h^6(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_6 + (\rho_a)_6 = -\mu_6(a_6) + (f_6)_6P_6) + \big((f_2)_6(Q_6)\big)(\rho_6)_6 \quad (4)$$

$$\rho_6(t_6, 0) = h_6(Q_6)\int_0^\infty \beta_6(a_6)\rho_6(t, a_6)da, \quad (5)$$

$$\frac{dQ_6}{dt}$$

$$= -b_6Q_6 + g_6(P_6)Q_6 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_6$ and $Q_6$ of (4-6) satisfy

$$0 \ \rho_6(t, a_6) \ \rho_6^*(t, a_6),$$

$$0 \leq Q_6(t) \leq Q_6^*(t),$$

Where $\rho_6^*(t, a_6)$ and $Q_6^*(t)$, are solutions of

$$(\rho_t^*)_6 + (\rho_a^*)_6 = -(\mu_6(a_6) + (f_6)_6(P_6^*) + (f_2)_6 Q_6^*))\,\rho_6^*, \quad (7)$$

$$\rho_6^*(t, 0) = \int_0^\infty \beta_6(a_6)\,\rho_6^*(t, a_6)\,da_6 \quad (8)$$

$$\frac{dQ_6^*}{dt} = -b_6\,Q_6^* + g_6(P_6^*)Q_6^*, \quad (9)$$

With $P_6^*(t) = \int_0^\infty \rho_6^*(t, a_6)da_6$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_6^*(t)}{dt} = P_6^*(t)\big(r_6 - (f_6)(P_6^*) - (f_2)_6(Q^*)\big), \quad (10)$$

$$\frac{dQ_6^*(t)}{dt} = Q_6^*(t)(-b_6 + g_6(P_6^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_6(a)e^{-\int_0^a \mu_6(\tau)dt}\,e^{-ra}\,da_6 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_6^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(t)dt}\, da_6 > 1, \qquad (H1)$$

Which guarantees $r_6 > 0$.

*Remark.* If $\beta_6(a) = [\beta_0]_6 a_6 e^{-\alpha_6 a_6}, \mu_6(a_6) \equiv \mu_6$ then the solution of (12) is

$$r_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_6]_6(P_6) = 0, [f_2](Q_6) = \lambda_6 Q_6\ h_6(Q_6) = 1, \qquad and\ \ g_6(P_6) = c_6 P_6$$

$$[\rho_6]_6 + [\rho_a]_6 = -(\mu_6 + \lambda_6 Q_6)\rho_6$$

$$\rho_6(t,0) = \int_0^\infty [\beta_0]_6 a_6 e^{-\alpha_6 a_6} \rho_6(t, a_6)\, da_6$$

$$\frac{dQ_6}{dt} = -b_6 Q_6 + c_6 P_6 Q_6$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_6}{dt} = \omega_6 P_6 - \lambda_6 P_6 Q_6$$

$$\frac{dQ_6}{dt} = b_6 Q_6 - c_6 P_6 Q_6$$

Where $\omega_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$

## 13. THE EQUILIBRIUM $(\rho_0)_6(a_6), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_6(a_6), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_6 + (\rho_a)_6 = -\big(\mu_6(a_6) + (f_6)_6(P_6)\big)\rho,$$

$$\rho_6(t,0) = \int_0^\infty \beta_6(a_6)\rho_6(t,a_6)da_6 \qquad (13)$$

The solution of (63) has the form of

$$\rho_6(t,a_6) = Ce_6^{r(t-a_6)} - \int_0^a (\tau)d\tau - \int_0^6 (f_6)_6\big(P_6(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_6(t) = C_6 e^{r_6 t - \int_0^t (f_6)_6(P_6(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_6^{-6}(r)$ if $k_6 > r$. This is equivalent to

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{(k_6)_6 a_6} da_6 < 1.$$

In fact, if we define

$$R_6(s_6) = \int_0^\infty \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau)d\tau} e^{-s_6 a_6}\, da_6$$

$R(s)$ is a decreasing function. $R_6(r_6) = 1$ And $R_6(k_6)_6 < 1$ imply $r_6 < (k_6)_6$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_6(a_6), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{k_6 a_6} da_6$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_6 \equiv \int_0^\infty (\rho_0)_{66}(a_6)da_6$ *satisfies*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{-(f_6)_6(P_0)a_6} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_6(a_6), 0)$

Let $u_6(t, a_6) = \rho(t, a_6) - \rho_0(a_6)$ Then

$$N_6(t) \equiv \int_0^\infty u_6(t, a_6)da_6 = \int_0^\infty \rho_6(t, a_6)da_6 - \int_0^\infty (\rho_0)_6(a_6)da = P_6(t) - P_0)_6$$

And

$$(u_6)_6 + (u_a)_6 = (\rho_t)_6 + (\rho_a)_6 - (\rho_0^6)_6(a_6)$$

$$= -\mu_6(a_6) + (f_6)_6(P_6) + (f_2)_6(Q_6))\,(u_6 + (\rho_0)_6) + \mu_6(a_6)(\rho_0)_6 + (f_6)_6(P_0)_6(\rho_0)_6$$

$$= -\mu_6(a_6)u_6 - (f_6)_6(P_6) + (f_2)_6(Q_6))\,u_6 - (\rho_0)_6\,(f_6)_6(P_6) - (f_6)_6(P_0)_6$$

$$+ (f_2)_6(Q_6)), \qquad (15)$$

Where $(\rho_0^6)_6(a_6)$ denotes $d(\rho_0)_6(a_6)/da_6$

Expansion $(f_6)_6(P_6)$ and $(f_2)_6(Q_6)$ at $(P_0)_6$ and 0 yield respectively

$$(f_6)_6(P_6) = (f_6)_6((P_0)_6 + N_6)$$

$$= (f_6)_6(P_0)_6 + (f_6^6)_6(P_0)_6 N_6 + o_6(N_6), \qquad (16)$$

$$(f_2)_6(Q_6)$$

$$= (f_2^6)_6(0)Q + o_6(Q_6) \quad (17)$$

Write $f_6(P_0) = f_{60}, f_6'(P_0) = f_{60}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_6)_6 + (u_a)_6 \approx -(\mu_6(a_6) + (f_{60})_6)u - (f_{60}^6)_6(\rho_0)_6(N_6) - (f_{20}^6)_6(\rho_0)_6(Q_6)$$

$$u_6(t,0) = \rho_6(t,0)(\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a) + (\rho_0)_6(a_6)da_6 - (\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a_6)da_6 + (h_6(Q_6) - 1)(\rho_0)_6(0) \quad (18)$$

Expanding

$$h_6(Q_6) = 1 + (h^6)_6(0)Q_6 + 0(Q_6) \equiv 1 + h_0'Q_6 + 0Q_6$$

And substituting it into (18), we arrive at

$$u_6(t,0) \approx \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + (h_{a_6}')(\rho_0)_6Q_6$$

The linearization is the following :

$$[u_6] + [u_6] = -(\mu_6(a_6) + [t_{60}]_6u_6 - [\alpha_6]_6(a_6)N_6 - [\alpha_2]_6(a_6)Q_6$$

$$u_6(t,0) = \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + [\xi_6]Q_6 \quad (19)$$

$$\frac{dQ_6}{dt} = -b_6Q_6$$

where

$$[\alpha_6]_6(a_6) = [t_6^6 a]_6[p, o]_6(a_6)d2(a_6) = [f_{20}^6]_6(\rho\, o\,)_6(a_6)\xi = [\mathrm{h\ o}^6]_6[\rho \mathrm{o}]_6(\mathrm{o}) \text{ and } \mathrm{b}_6 = \mathrm{b}_6 - g_6(P_6)$$

If $g_6(P_0) > b$, clearly, $(\rho_0)_6(a_6), 0)$ is unstable and the predator population is persistent. Assume $g_6(P_0)_6 < b$.

Then, $Q_6(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_6 + u_a = -(\mu_6(a_6) + (f_{60})_6)u_6 - (\alpha_6)_6(a_6)N_6$$

$$= -(\mu_6(a_6) + f_{60})u_6 - \alpha_6(a_6)\int_0^\infty u_6(t, a_6)da_6$$

$$u_6(t, 0) = \int_0^\infty \beta_6\,(a_6)u_6(t, a_6)da_6$$

Using the method developed by Cushing [2, 3], we suppose

$$u_6(t, a_6) = \tilde{\mathrm{u}}_6(a_6)^{e^{\gamma(t-a_6)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_6}\, da_6$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_6)$ gives,

$$\tilde{\mathrm{u}}_6(a_6) = e^{-\int_0^a \mu_6(\tau_6)d\tau - (f_{60})_6 a}\left(\tilde{\mathrm{u}}_6(0) - w\int_0^a e^{\int_0^\tau \mu_6(s_6)ds + (f_{60})}\tau_6\,(\alpha_6)_6(\tau_6)e^{\gamma_6\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_6(a_6)e^{-\gamma_6 a_6 - (f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6)d\tau}\right)\tilde{\mathrm{u}}_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a(\alpha_6)_6(\tau_6)e^{-\gamma_6\tau_6}e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds}d\tau_6\,da_6$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}da\right)(\tilde{u}_6(0))_6$$

$$+\left(-6-\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds+y\tau}\alpha_6(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(t_6)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_6\,(a_6)Y_6(\gamma_6,a_6)da_6\right)u_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6=0,$$

(24)

$$\left(\int_0^{\infty}Y_6(\gamma_6,a_6)da_6\right)\tilde{u}_6(0)$$

$$+\left(-1-\int_0^{\infty}Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6\right)w=0.$$

Equqtion (20) has a solution of the form of $u_6(t,a_6)=u_6(a_6)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_6(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) Y^{-6}(\gamma_6, a_6) d\tau_6$. We obtain the characteristic equation

$$f_{60}^{l}{}_6 \left( \int_0^\infty Y_6(\gamma_6, a_6) da_6 \right) \left( \int_0^\infty \beta_6(a_6) Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$+ \left( 1 - \int_0^\infty \beta_6(a_6) Y_6(\gamma_6, a_6) da \right) \left( 1 + (f_{60}^{l})_6 \int_0^\infty Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$= 0. \qquad (25)$$

Since

$$(\rho_0)_6(a_6) = (\rho_0)_6(0) e^{-\int_0^a \mu_6(\tau_6) d\tau_6 - (f_{60})_6 a_6}$$

$$I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) e^{\gamma_6 \tau_6 + (f_{60})_6 \tau + \int_0^\tau \mu_6(s_6) ds_6} d\tau$$

$$= \frac{(\rho_0)_6(0)}{\gamma_6} (e^{\gamma_6 a_6} - 1)$$

Let

$$E_6(a_6) = e^{-(f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6) d\tau_6}$$

This relationship is cardinal one. It is useful in quantum computation and other black hole dynamics problems

**Quantum space and quantum time**

(632)**Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.**

(633) **In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.**

(634) **They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.**

(635)**When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.**

(636) **Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.**

**The Planck scale**

(637)**Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".**

(638)**Here ħ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.**

(639)**The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.**

(640)**Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)**

(641)**Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.**

## Physical applications

**Some phenomena where quantum gravity plays a major role are the following.**

**The microscopic structure of spacetime**

(642)**If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.**

(643) **Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length**

(644) **Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.**

(645) **A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.**

**Early cosmology**

(646)**According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.**

(647) **However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.**

(648) **A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.**

(649) **Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".**

(650)**One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.**

**Black holes**

(651)**Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.**

(652) **First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.**

(653) **Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.**

(654)**What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.**

(655) **Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.**

**Astrophysical effects**

(656)**Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.**

(657) **An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.**

(658)**The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.**

**Main approaches**

(659)**The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.**

(660) **Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.**

(661)**The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.**

**String theory**

(662)**String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.**

(663) **Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.**

(664)**String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.**

(665)**Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.**

(666)**However, when summed the entire series appears to be (=) divergent.**

(667)**In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.**

(668)**Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.**

(669)**Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.**

(670) **In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.**

**Loop quantum gravity and spin foams**

(671)**Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.**

| |
|---|
| (672)**The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity.** |
| (673)**Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime.** |
| (674)**The granular properties of space can be (=) explicitly computed.** |
| (675) **In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory.** |
| (676)**The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks.** |
| (677) **Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure.** |
| (678) **The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit.** |
| (679)**Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism.** |
| (680)**This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism.** |

**Noncommutative geometry**

| |
|---|
| (681)**Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime.** |
| (682) **Quantum mechanics teaches us that dynamical quantities are (=) noncommutative --meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed.** |
| (683)**It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry.** |
| (684) **A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces.** |
| (685)**Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces.** |
| (686)**This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity.** |

**Others**

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|---|
| (687) **Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays** |

(688)**. Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection.**

(689) **Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible.**

(690) **The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.**

**Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016**

(691)**These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field.**

(692)**The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets.**

(693) **Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state.**

(694) **Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands.**

(695)**The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI**

(696) **The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9].**

(697)**This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive.**

(698)**More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.**

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

(699)**Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.**

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| (700)**Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.** |
| (701) **The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.** |
| (702)**This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.** |
| (703)**The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.** |
| (704)**This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.** |
| (705) **For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.** |
| (706)**The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society** |

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

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| (707)**Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.** |
| (708)**We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.** |
| (709)**When this is true—as is often the case—the system serves as (=) a "bath" for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called thermalization.** |
| (710)**However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.** |
| (711)**Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1).** |
| (712)**Bringing them to thermal equilibrium requires (e) an external bath.** |
| (713) **But so far, there has been no rigorous proof of (e) the existence of such many-body localization. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption.** |
| (714)**The concept of localization was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations).** |

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| (715)**The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century.** |
| (716)**But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information.** |
| (717)**Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment.** |
| (718)**For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood.** |
| (719) **It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7].** |
| (720) **Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature.** |
| (721) **One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.** |
| (722)**However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].** |
| (723)**Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.** |
| (724) **Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.** |
| (725)**As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked.** |
| (726)**His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice.** |
| (727) **An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site.** |
| (728)**This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself).** |
| (729)**This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition.** |
| (730)**However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime:** |

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| (731) **These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects.** |
| (732)**It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies).** |
| (733) **This is a realistic scenario for (e) all known systems.** |
| (734)**On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects.** |
| (735)**On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other.** |
| (736)**No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true.** |
| (737) **His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found.** |
| (738)**Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization?** |
| (739)**Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects?** |
| (740)**Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions.** |
| (741)**So in those systems, the quantum phase transition does indeed exist.** |
| (742)**But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension.** |
| (743)**Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters.** |

**Focus: Biological Cells Form Electric Circuits July 1, 2016• Physics 9, 75 Cells that are electrically active and that also produce light for easy voltage monitoring could lead to new studies of heart arrhythmias and possibly bio-computing. Figure caption expand figure H. M. McNamara et al., Phys. Rev. X (2016) Cell circuit. Waves of electrical activity propagate in a loop within a collection of cells engineered with electrical and optical properties (see video below).Some sentences, pictures and videos are deleted due to constraints on space.**

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| (744)**The human heartbeat is produced by (e) electrical pulses that propagate through (e&eb) cardiac tissue, causing (eb) rhythmic muscle contraction.** |
| (745)**Researchers have previously engineered cells to create (eb) an artificial tissue capable of producing (eb) coordinated electrical activity, and now a team has added the ability to monitor (e&eb) their electrical state by (e) detecting fluorescent emission.** |
| (746)**They have also fashioned the cells into (e&eb) "living circuits" that might act as (=) model systems for studying heart behavior.** |
| (747)**Constructing electrically active tissues from their basic building blocks by (e) genetically engineering cells allows (eb) researchers to study a "clean" system similar to (e) the more complicated cardiac tissue.** |

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| (748)**These systems let them investigate, for example, the electrical signals that lead to (eb) the potentially fatal irregular heartbeat called arrhythmia.** |
| (749)**Previously researchers have introduced (eb) ion channels—proteins that conduct ions through (e&eb) cell membranes—into human kidney cells to make (eb) electrically "excitable" tissues [1].** |
| (750)**Adam E. Cohen of Harvard University and his co-workers decided to add more ingredients to such cells, so that the electrical activity could be (=) light-activated and easily visualized.** |
| (751)**They incorporated four additional gene products into (e&eb) human embryonic kidney (HEK) cells.** |
| (752) **Two were ordinary ion channels; the other two were (=) a blue-light-activated ion channel (a light-sensitive switch) and (e&eb) a protein that fluoresces with (e&eb) an intensity that depends on (=) the voltage across the membrane in which it is embedded.** |
| (753)**Light activation of electrically excitable cells has been shown before [2]. But Cohen and colleagues have now included a convenient "readout" mechanism (the fluorescent protein) and have used patterning methods to shape (e&eb) the cell cultures into "tissue circuits."** |
| (754)**They grew rings of cells a few millimeters in diameter that could act as (e&eb) oscillators that support (eb) circulating waves of electrical excitation.** |
| (755)**These waves arise from (e) the collective behavior of the cells: a pulse produced by (e) one cell can trigger another from (e) its neighbors, eventually producing (eb) a coherent wave of excitation traveling through (e&eb) the tissue.** |
| (756) **It's this coordination, in a repeating cycle of excitation and quiescence, which gives rise to (eb) rhythmic activity in cardiac tissue.** |

**H. M. McNamara et al., Phys. Rev. X (2016)**

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| (757)**A pattern of blue light (not visible) illuminates these cells, defining (eb) a rectangular loop around which electrical waves can propagate.** |
| (758)**At about 0:07 the rectangle begins to shrink in (eb) the vertical direction, causing t(eb) he cycle time to gradually decrease. (In other experiments, the cells were grown in (eb) fixed, ring-shaped patterns.)** |
| (759)**The ring of engineered HEK cells acted like an electronic "ring oscillator," a device with a characteristic frequency to (e&eb) its circulating electrical signal.** |
| (760)**The electrical waves in a ring oscillator can propagate (e&eb) in either direction around the ring, enabling (eb) it to serve as a binary memory unit.** |
| (761) **Cohen and colleagues found that in their HEK tissues, these circulating waves remained (eb) stable for more than two hours, with (e&eb) the pulses completing a few circuits every second.** |
| (762)**The electrical signals were easily visualized as (=) waves of fluorescent emission.** |
| (763)**The researchers say that these tissues could be used to (e) understand and predict the behavior of real cardiac cells, especially arrhythmias.** |
| (764)**"In real heart, if you pace the tissue fast enough, *(eb) it starts skipping beats and eventually becomes (=) arrhythmic," says Cohen.** |
| (765) **This state is a well-known precursor to (e) sudden cardiac death and is associated with (e&eb) the breakup of the usual regular waves of electrical excitation.** |
| (766) **Some drugs also have the dangerous side effect of inducing (eb) arrhythmia, and Cohen says that the new system might offer a quick and convenient way of screening for such hazards. The work "is really very elegant and deserves attention," says physiologist Gil Bub of Oxford University in the** |

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| **UK. "It opens new research questions for people like me who are interested in wave dynamics in excitable systems."** |
| (767)**There might also be practical uses for information processing in living tissues. "By combining patterns of cell growth and (e&eb) patterns of ion channel expression, we think we can build diodes, amplifiers, and other simple logic elements," says Cohen.** |
| (768)**Andrew Adamatzky of the University of the West of England in Bristol, a specialist in "unconventional computing," foresees circuits that perform computations through (e&eb) the interactions of electrical waves. "This research might lead to fabrication of implantable bio-computers," he says. This research is published in Physical Review X. –Philip Ball Philip Ball is a freelance science writer in London and author of Invisible: The Dangerous Allure of the Unseen (University of Chicago Press, 2014).** |

**Optically Controlled Oscillators in an Engineered Bioelectric Tissue Harold M. McNamara, Hongkang Zhang, Christopher A. Werley, and Adam E. Cohen Phys. Rev. X 6, 031001 – Published 1 July 2016**

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| (769)**Complex electrical dynamics in excitable tissues occur throughout (e&eb) biology, but the roles of individual ion channels can be difficult to determine due to (e) the complex nonlinear interactions in (eb) native tissue.** |
| (770) **Here, authors ask whether they can engineer a tissue capable of (e) basic information storage and processing, where (e) all functional components are known and well understood.** |
| (771)**Authors develop a cell line with (e&eb) four transgenic components: two to enable (eb) collective propagation of electrical waves and (e&eb) two to enable optical perturbation and (e&eb) optical readout of membrane potential.** |
| (772)**They pattern the cell growth to define (eb) simple cellular ring oscillators that run stably for >2 h (∼104 cycles) and that can store data encoded in (eb) the direction of electrical circulation.** |
| (773) **Using (e) patterned optogenetic stimulation authors probe the biophysical attributes of (e) this synthetic excitable tissue in detail, including (e) dispersion relations, (e&eb) curvature-dependent wave front propagation, (e&eb) electrotonic coupling, and (e&eb) boundary effects.** |
| (774)**Authors then apply the biophysical characterization to (e) develop an optically reconfigurable bioelectric oscillator.** |
| (775)**These results demonstrate the feasibility of engineering bioelectric tissues capable of (e&eb) complex information processing with optical input and output. Received 12 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevX.6.031001 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society** |

**Scientists detect radio emission from a nearby brown dwarf July 6, 2016 by Tomasz Nowakowski report**

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| (776)**K2 and Spitzer images of W0607+24: The position of W0607+24 is marked by (e) a red circle. SDSS J060738.06+242939.5, marked by (e) an orange circle, are the brightest star in the K2 image.** |
| (777)**The three faint (r = 18.3 − 19.6) stars marked by cyan circles, contribute significant flux to (e) the K2 measurement but are (-) negligible at Spitzer's longer wavelength and (e &êb) higher resolution. Other stars with r < 20 in SDSS are marked with blue circles. Credit: Gizis et al., 2016.** |

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| (778)**(Phys.org)—Located some 23 light years away, the brown dwarf designated WISEP J060738.65+242953.4 or W0607+24 for short, turns out to be (=) a source of radio emission.** |
| (779)**According to a research paper published July 4, this substellar object showcases (eb) quiescent radio emission, making it one of (e) the most radio-faint, ultra-cool dwarfs yet detected.** |
| (780)**Brown dwarfs like W0607+24 are objects that are too large to be called planets and too small to be stars. With a mass below that necessary to maintain hydrogen-burning nuclear fusion reactions, they are much cooler and dimmer than (eb) main sequence stars.** |
| (781)**These objects could be also sources of (e) radio emission, but the cause of this process is still not completely understood.** |
| (782)**Due to its proximity, W0607+24 is (=) an excellent target for studying the nature of (e) brown dwarfs.** |
| (783) **It is the nearest known late-L dwarf in the northern hemisphere, and the third-nearest in (=) the whole sky, (eb) allowing researchers to study phenomena occurring in brown dwarfs.** |
| (784)**That's why a team of astronomers, led by John Gizis of the University of (e)Delaware, Newark, carried out a series of mid-infrared, radio, and optical observations of W0607+24 to study any activity in this curious substellar object.** |
| (785)**Motivated by the results from NASA's prolonged Kepler mission, named K2, which obtained (eb) long time-series photometry of W0607+24; they observed this object simultaneously with the Karl G. Jansky Very Large Array (VLA) in New Mexico and NASA's Spitzer Space Telescope.** |
| (786)**The most important results were provided by VLA, as the team found that W0607+24 is (=) a source of quiescent, non-bursting, radio emission.** |
| (787)**The scientists noted that the discovery of this process in this nearby object could be (=) crucial for our understanding of this process in (eb) ultracool dwarfs in general.** |
| (788)**"W0607+24 is a valuable test case for (e) understanding the mechanisms of radio emission in ultracool dwarfs.** |
| (789)**Some of the radio-detected objects emit (eb) rapid, intense bursts with high circular polarization, often found to occur periodically and attributed to coherent emission due to (e) the electron cyclotron maser instability," the paper reads.** |
| (790)**The astronomers also suggested that the relative radio faintness of W0607+24 may be due to (e) its orientation if (e) the emission is concentrated at the poles.** |
| (791)**However, follow-up observations are needed to better characterize (eb) the detected radio emission and to eliminate (e) all the uncertainties.** |
| (792)**The team also hopes for more significant findings when (e) studying this nearby substellar object in the future.** |
| (793)**"As the nearest known northern hemisphere late-L dwarf, this offers many possibilities for future studies. For example, it could serve as (=) an useful standard for (e) line-broadening and (e&eb) polarization studies.** |
| (794)**If there is a planetary system aligned with (e&eb) (the primary's rotation axis, then we would expect very little radial velocity signal but a strong astrometric signal," the researchers concluded.** |
| (795)**Moreover, the team plans to monitor (e&eb) additional late-M and L dwarfs in other K2 campaigns, expecting even more promising results.** |
| (796) **They are more distant than W0607+24, but because they are warmer, (eb) some are brighter and higher signal-to-noise than that described in this paper.** |

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| (797)**In addition to detecting the radio emission, (eb)the scientists also estimated the age, mass and radius of W0607+24** |
| (798)**. According to the study, the object is less than two billion years old, has (e) a maximum mass of about 0.055 solar masses and a radius approximately of 0.1 solar radii.** |
| (799)**More information: WISEP J060738.65+242953.4: A Nearby. Pole-On L8 Brown Dwarf with Radio Emission, arXiv: 1607.00943 [astro-ph.SR] arxiv.org/abs/1607.00943** |
| (800)**Authors present a simultaneous, multi-wavelength campaign targeting the nearby (7.2 pc) L8/L9 (optical/near-infrared) dwarf WISEP J060738.65+242953.4 in (eb) the mid-infrared, radio, and optical.** |
| (801)**Spitzer Space Telescope observations show no (e) variability at the 0.2% level over 10 hours each in the 3.6 and 4.5 micron bands.** |
| (802) **Kepler K2 monitoring over 36 days in Campaign 0 rules out (e) stable periodic signals in (eb) the optical with amplitudes great than 1.5% and periods between 1.5 hours and 2 days.** |
| (803)**Non-simultaneous Gemini optical spectroscopy detects lithium, constraining (e) this L dwarf to be less than ~2 Gyr old, but no Balmer emission is observed.** |
| (804)**The low measured projected rotation velocity (v sin i < 6 km/s) and lack of variability are (=) very unusual compared to other brown dwarfs, and authors argue that this substellar object is (=) likely viewed pole-on.** |
| (805)**Authors detect quiescent (non-bursting) radio emission with (e&eb) the VLA.** |
| (806) **Amongst radio detected L and T dwarfs, it has (e) the lowest observed L_nu and the lowest v sin i.** |
| (807)**They discuss the implications of pole-on detection for (e) various proposed radio emission scenarios. © 2016 Phys.org** |

**A new look at the galaxy-shaping power of black holes July 6, 2016**

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| (808)**It wasn't until the advent of X-ray astronomy that the full picture of (e) the galaxy-shaping power of black holes began to emerge (eb) with the ability to see plasma.** |
| (809)**Data from a now-defunct X-ray satellite is providing (eb) new insights into the complex tug-of-war between galaxies, the hot plasma that surrounds them, and the giant black holes that lurk in their centres.** |
| (810)**Launched from Japan on February 17, 2016, the Japanese space agency (JAXA) Hitomi X-ray Observatory functioned for just over a month before contact was lost and (e&eb) the craft disintegrated.** |
| (811)**But the data obtained during those few weeks was enough to paint (eb) a startling new picture of the dynamic forces at work within galaxies.** |
| (812)**New research, published in the journal Nature today, reveals data that shows just how important the giant black holes in galactic centres are to the evolution (eb) of the galaxies as a whole.** |
| (813)**"We think that supermassive black holes act(e&eb) like thermostats," said Brian McNamara, University Research Chair in Astrophysics at the University of Waterloo. "They regulate (e&eb) the growth of galaxies."** |

**Champagne bubbles of plasma**

(814)**During its brief life, the Hitomi satellite collected X-ray data from (e) the core of the Perseus cluster, an enormous gravitationally-bound grouping of (e) hundreds of galaxies.**

(815) **Located some 240 million light years from earth, the Perseus cluster is one of (e) the largest known structures in the universe.**

(816) **The cluster includes not only the ordinary matter that makes up the galaxies, but an (e) "atmosphere" of hot plasma with a temperature of tens of millions of degrees, as well as a halo of (e) invisible dark matter.**

(817)**Earlier studies, going back to the 1960s, have shown that each of the galaxies in the cluster - and indeed most galaxies - likely contains (e) a supermassive black hole in its centre, an object 100 million to more than ten billion times as massive as our sun**

**New look at the galaxy-shaping power of black holes**

(818)**Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving.**

(819)**"These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science.**

(820) **"Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy."**

(821)**This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return.**

(822)**The remaining matter gets absorbed into the black hole, adding to (e) its mass.**

(823)**The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne.**

(824)**The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution.**

(825)**Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form.**

(826)**"Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara.**

(827)**"It creates (eb) these bubbles, and the bubbles keep the plasma hot.**

(828)**That's what prevents galaxies from (e) becoming even bigger than they are now."**

(829)**Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.**

(830)**In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.**

(831)**Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.**

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| (832)**"It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.** |
| (833)**"It's like nature's thermostat, which keeps these galaxies from (e) growing.** |
| (834) **If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."** |

**A new look at the galaxy-shaping power of black holes**

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| (835)**An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the** |
| (836)**McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.** |
| (837)**Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.** |

**Doomed satellite missions**

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| (838)**Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.** |
| (839)**At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.** |
| (840)**Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.** |
| (841)**The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.** |
| (842)**These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.** |
| (843)**"But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people** |
| (844)**. It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."** |
| (845)**The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.** |
| (846) **"This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More** |

**information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website**

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(847)**Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.**

(848)**They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.**

(849)**However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.**

(850) **Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.**

(851)**X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well.**

(852)**The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma.**

(853)**These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6.**

(854)**Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus.**

(855)**A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate.**

(856)**Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure.**

**Dark matter guides growth of supermassive black holes February 18, 2015**

(857)**This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy.**

(858)**Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole.**

(859) **But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy.**

(860)**A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole.**

(861) **It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth.**

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| (862)**"There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA).** |
| (863)**This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies.** |
| (864)**However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated.** |
| (865)**In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1.** |
| (866) **We know dark matter exists only from (e) its gravitational effects.** |
| (867) **It holds together galaxies and (e&eb) galaxy clusters.** |
| (868) **Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years.** |
| (869)**To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies.** |
| (870) **They used star motions as tracers to (e) weigh the galaxies' central black holes.** |
| (871) **X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto.** |
| (872)**They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone.** |
| (873)**This connection is likely to be related to (e&eb) how elliptical galaxies grow.** |
| (874) **An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together.** |
| (875)**Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole.** |
| (876)**"In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp** |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (877)**The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the …** |
| (878)**Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness?** |

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| (879)**Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness.** |
| (880)**"Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain."** |
| (881)**Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness.** |
| (882)**"With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication."** |
| (883)**Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories.** |
| (884)**The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen.** |
| (885) **In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer.** |
| (886)**Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants.** |
| (887)**These were treated as (=) "aware" and "unaware" trials, respectively.** |
| (888)**Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it.** |
| (889)**The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another.** |
| (890)**Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion.** |
| (891) **No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness.** |

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| (892)**"We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said.** |
| (893)**"Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness."** |
| (894)**The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner.** |
| (895)**Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain.** |
| (896)**"We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said.** |
| (897) **"This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients** |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

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| (898)**Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter.** |
| (899) **As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined.** |
| (900)**Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions.** |
| (901)**"Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.** |
| (902)**The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.** |
| (903)**However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.** |
| (904)**The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.** |
| (905) **Noda and his colleagues instead studied the change in emissions over (e&eb) time.** |

(906) **By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.**

(907)**Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).**

(908)**The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.**

(909)**Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.**

(910)**"The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp**

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(911)**The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking "palm tree" organization.**

(912)**The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society**

(913)**(Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible?**

(914) **With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem.**

(915)**This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other.**

(916)**According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible.**

(917)**This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges.**

(918) **In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century.**

(919)**Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching.**

(920)**In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom.**

(921) **Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth.**

**Spheres on trees**

(922)**Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration.**

(923)**They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement.**

(924) **Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum.**

(925)**By visualizing the problem in this way, the researchers noticed that these particular graphs don't (e) have lots of branches extending from lots of (e) other branches.**

(926) **Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure.**

(927)**The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network.**

(928)**As a result, it provides (eb) an important clue for finding the global minimum.**

(929)**It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum.**

(930)**This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large.**

**Thompson's problem 2**

(931)**By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society**

(932)**The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima.**

(933) **The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node.**

**Implications for Thomson's and Smale's problems**

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| (934)**Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before.** |
| (935)**"Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around."** |
| (936)**As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem.** |
| (937)**"I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view."** |
| (938)**In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges.** |
| (939) **From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear.** |
| (940)**"While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK.** |
| (941)**"For larger numbers, there will likely be (=) multiple funnels.** |
| (942)**Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum."** |
| (943)**Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.** |
| (944)**"The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.** |
| (945)**"Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org** |

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**Introduction:**

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price,**

**University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss.**

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| **(97)** | **Black Holes are regions of (e) space in which gravitational fields are so strong that no particle or signal can escape (eb) the pull of gravity.** |
| **(98)** | **The boundary of this no-escape region is (=) called the event horizon, since distant observers outside the black hole cannot (e) see (cannot get light from) events inside.** |
| **(99)** | **Although the fundamental possibility of such an object exists within (eb) Newton's classical theory of gravitation, Einstein's theory of gravity makes (eb)black holes inevitable under some circumstances.** |
| **(100)** | **Prior to the early 1960s, black holes seemed to be (=) only an interesting theoretical concept with no (e) astrophysical plausibility, but with (e&eb) the discovery of quasars in 1963 it became clear that very exotic astrophysical objects could exist.** |
| **(101)** | **Nowadays it is taken for granted that black holes do exist in (eb) at least two different forms.** |
| **(102)** | **Stellar mass black holes are (=) the endpoint of the death of some stars, and supermassive black holes are the result of (e) coalescences in the centers of most galaxies, including our own.** |
| **(103)** | **No signal can propagate from (e) inside a black hole, but the gravitational influence of a black hole is (=) always present.** |
| **(104)** | **This influence does not (e) propagate out of the hole; it is permanently present outside, and depends only on (=) the total amount of mass, angular momentum, and electric charge that have gone into forming (eb) the hole.** |
| **(105)** | **Black holes can be detected through the influence of this (e&eb) strong gravity on the surroundings just outside the hole.** |
| **(106)** | **In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them.** |
| **(107)** | **Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors.** |

## Classical vs. relativistic black holes

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| **(108)** | **Something like a black hole exists within (eb) Newton's classical theory of gravity.** |
| **(109)** | **In that theory, an energy argument tells us that there is (=) an escape velocity $v_{esc}=\sqrt{2GM/R}$ from (e) the surface of any spherical object of (e) mass M and radius R** |
| **(110)** | **If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity.** |
| **(111)** | **Thus the condition for such an "unseeable" object is $R<2GM/c^2$** |
| **(112)** | **In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape.** |
| **(113)** | **This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity.** |
| **(114)** | **To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it.** |

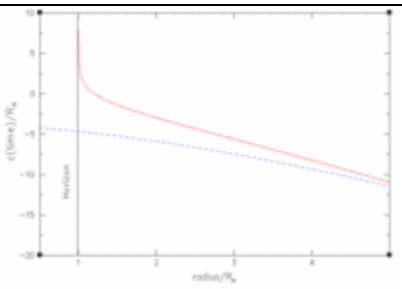
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| **(115) The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c .** |
| **(116) This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed.** |
| **(117) An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity.** |
| **(118) Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable.** |

**Stationary black holes**

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| **(119) In the Newtonian theory, gravity is described by (e) the potential Φ.** |
| **(120) Inside a spherical object the form of (e) Φ(r) depends on (=) the interior structure, but in the vacuum outside (e) matter the potential −GM/r depends only on (=) the interior mass.** |
| **(121) Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object.** |
| **(122) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward.** |
| **(123) This horizon radius RH=2GM/c2 is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory.** |
| **(124) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution.** |
| **(125) Radius RHhere actually means that the area of the event horizon is (=) 4πR2H.** |
| **(126) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution.** |
| **(127) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity.** |
| **(128) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own.** |
| **(129) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole.** |
| **(130) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density.** |
| **(131) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed.** |
| **(132) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole.** |
| **(133) In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter.** |

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| **(134) Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum.** |
| **(135) In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior.** |
| **(136) The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object.** |
| **(137) Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum.** |
| **(138) These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.]** |
| **(139) All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time.** |
| **(140) But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?"** |
| **(141) The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point.** |
| **(142) One of these coordinates is (=) called the "coordinate time."** |
| **(143) Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate.** |



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| **(144) Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole** |
| **(145) At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole.** |
| **(146) Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations.** |
| **(147) For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree** |
| **(148) An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events.** |

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| **(149) The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving.** |
| **(150) Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle.** |
| **(151) The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return.** |
| **(152) But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon.** |
| **(153) Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future.** |
| **(154) The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?"** |

**Black hole parameters**

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| **(155) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin).** |
| **(156) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun.** |
| **(157) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits.** |
| **(158) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form.** |
| **(159) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds.** |
| **(160) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2.** |
| **(161) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed.** |
| **(162) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer.** |
| **(163) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2)** |
| **(164) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials.** |
| **(165) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable** |
| **(166) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite.** |

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| **(167) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole.** |
| **(168) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars.** |
| **(169) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity.** |
| **(170) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity.** |
| **(171) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite.** |
| **(172) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity.** |
| **(173) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum.** |
| **(174) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship.** |
| **(175) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= (4πG2/c4)[(M+M2−a2−−−−−−−√)2+a2].** |
| **(176) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored.** |

**Dynamical black holes**

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| **(177) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape.** |
| **(178) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, a horizon can be born.** |
| **(179) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid.** |
| **(180) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward.** |
| **(181) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations.** |
| **(182) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes.** |
| **(183) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3.** |

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| **(184) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/** |
| **(185) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers.** |
| **(186) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission.** |
| **(187) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.** |
| **(188) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.** |
| **(189) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.** |
| **(190) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.** |
| **(191) The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature 6×10−8 (Mo/M) K.** |
| **(192) Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.** |

Then $[Y_6]_6[\gamma_6]_6, a_6) = E_6(a_6)e^{-\gamma_6 a_6}$. by a simple calculation,

$$\int_0^\infty \beta_6(a_6)Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma}\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}\ da\right);$$

$$\int_0^\infty Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma_6}\left(\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}\ da_6\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_6(0)}{\gamma_6}[f_{60}^6]\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}da_6\ (\int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6)$$

$$+\left(1-\int_0^{\infty}\beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right).$$

$$\left(1+\frac{(\rho_0)_6(0)}{\gamma_6}(f_{60}^{'})_6\left(\int_0^{\infty}E_6(a_6)da_6-\int_0^{\infty}E_6(a_6)e^{-\gamma_6 a_6}da_6\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right)(1+\frac{(f_{60}^{'})_6}{\gamma_6}(\rho_0)_6(0)\int_0^{\infty}E_6(a_6)da_6)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_6(a_6)Y_6(\gamma_6 a_6)da_6=1-\int_0^{\infty}\beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\neq 0,$$

Since $\int_0^{\infty}\beta_6(a_6)Y_6(\gamma_6 a_6)\int_0^{a}\alpha_6(\tau_6)Y^{-6}(\gamma_6,\tau_6)d\tau_6 da_6>0$. Therefore, the characteristic solution is

$$\gamma_6=-(f_{60}^{'})_6(\rho_0)_6\int_0^{\infty}E_6(a_6)da_6.$$

Since $f_6^{'}(P_0)_6>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_6),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE OF DEBIT TYPES AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_6), 0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_6 + (\mu_6(a_6) + [f_6]_6(P_6^*) + [f_2]_6(Q_6^*))P_6^* = 0, \quad (26)$$

$$p_6^*(0) = h_6(Q_6^*)\int_0^\infty \beta_6(a_6)p_6^*(a_J)p_6^*(a_6)da_6 \quad (27)$$

$$g_6(P_6^*) - b_6 = 0, \quad (28)$$

Where $P_6^* = \int_0^\infty p_6^*(a_6)\, da_6$

Equation (28) yields $P^* = g^{-6}(b)$,

From (26),

$$p_6^*(a) = p_6^*(0)e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6}$$

Substituting into (27), we need to solve

$$\left(1 - h_6(Q_6^*)\int_0^\infty \beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6}da_6\right)p_6^*(0) = 0$$

For $Q_6^*$

Fix $p_6^*(0)$ and $P_6^*$ Let

$$H_6(z_6) = 1 - h_6(z_6)\int_0^\infty \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6}\, da_{,6}$$

And assume that

$$[k_3]_6 \int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(g_6^{-6}(b_6))a_6 - [k_2]_6 a_6}\, da_6 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_6(z_6)] > 0.$$

Clearly,

$$\frac{dH_6}{dz_6} = h_6(z)([f_2^{'}]_6(z_6))\int_0^\infty a_6\beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - f_6(P_6^*)a_6 - f_2(z_6)a_6}\, da_{6}$$

$-h^{'}(z_6)\int_0^\infty \beta_6(a_6)\, e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6}\quad da_6$

By the assumptions on $\left[f_2^6\right]_6 and\ [h^6]^6 dH_6/dz > 0$. Hence, if

$$H_6(0) = 1 - \int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6} da_6 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_6(0) < 0$ follows if $g_6^-(b_6) < [f_6^{-6}]_6(r_6)$, where $r_6$ is the solution of (12).

Indeed, since

$$R_6(s_6) = \int_0^{\infty} \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 e^{-y_6^a} da_6}$$

Is a decreasing function of $s_6$ when $P_6^* = [g^{-6}]_6(b_6) < [f_6^{-6}]_6(r_6)$, that is $[f_6]_6(P_6^*) < r$, $R\big(f_6(P^*)\big) > 1$which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_6(g^{-6}(b))) > 6, i.e.,$

$\int_0^{\infty} \beta_6(a_6) e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6\left(g_6^{-6}(b)\right)a_6} > 1.$ (H4)

Then the system (4)-(6) has a unique positive equilibrium $(P_6^*(a_6), Q_6^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_6^*(a_6), Q_6^*)$ is more complicated. We again consider perturbations $m_6(t_6 a_6) = p_6(t_6, a_6) - p_6^*(a_6)$ and $W_6(t) = Q_6(t) - Q^*$.Then

$M_6(t) \equiv \int_0^{\infty} m_6(t, a_6) da_6 = P_6(t) - P_6^*,$

$$[m_6]_6 + [m_a]_6 = [p_6]_6 + [p_a]_6 - \left[p^*\right]_6(a_6) \approx -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6$$

$$-[f_{6*}^6]_6 p_6^*(a_6)M - [f_{2*}^6]_6 p_6^*(\mathrm{a})W_6,$$

$$m_6(t,0)p_6(t,0) - p_6^*(0) \approx h_* \int_0^{\infty} \beta_6(a_6) m_6(t, a_6)\, da + \frac{[h_*^6]_6}{[h_*]_6}\, p_6^*(0)W_6$$

$$\frac{dW_6}{dt} = \frac{dQ_6}{dt} \approx [g_*^6]_6 Q_6^* M_{,6}$$

Where $[f_{6*}]_6$ denoting $[f_6(P^*)]_6$, $[f_{6*}]_6$ denoting $[f_2(P^*)]_6$, $[f_{6*}']_6$ denoting $[f'_6(P^*)]_6 [f_{2*}']_6$ denoting $[f_2]_6$ $(Q_6^*,\quad h_6^*$ denoting $h_*'$ denoting $h_6'$ $(Q_6^*)$, and $[g_*']_6$

Denoting $[g']_6(P_6^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_6]_6 + [m_a]_6 = -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6 - [f'_{6*}]_6[p_6^*](a_6)M_6 - [f_{2*}']_6 p_6^*(a_6)W_6$$

$$m_6(t,0) = [h_*]_6 \int_0^\infty \beta_6(a_6)m_6(t,a_6)da_6 + \frac{[h_*^6]_6}{[h^*]_6}[p^*]_6(0)\, W_{,6} \qquad (30)$$

$$\frac{dW_6}{dt} = [g_*^6]_6[Q_6^*]M_6$$

Integrating $[m_6]_6(t,a_6)$ with respect to a from 0 to ∞, we have

$$\frac{dM_6(t)}{dt} = \int_0^\infty [mt]_6(t,a_6)da_6$$

$$= -\int_0^\infty [m_a]_6\,(t,a_6)da_6 - \int_0^\infty \mu(a_6)m(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + [f'_{6*}]_6 \int_0^\infty p_6^*(a_6)da_6\right)M_6 - [f_{2*}^6]_6 \int_0^\infty p_6^*(a_6)da_6 W_6$$

$$= m_6(t,0) - \int_0^\infty \mu_6(a_6)\, m_6(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + [f_{6*}']_6 P_6^*\right)M_6 - [f_{2*}']_6 P_6^* W_{,6}$$

From (30), we arrive at the equations

$$\frac{dW_6(t)}{dt} = \int_0^{\infty} ([h_*]_6\beta_6(a_6) - \mu_6(a_6))m_6(t,a_6)da_6$$

$$-\left([f_{6*}] + [f_{2*}]_6 + \left[f_{6*}^{'}\right]_6 P_6^*\right)M_6 + \left(\frac{[h_*^{'}]_6}{[h_*]^6}\ P_6^*(0) - f_{2*}^{'}P_6^*\right)W_6,$$

(36)

$$\frac{dW_6(t)}{dt}\left[g^{'}{}_*\right]_6 Q_6^*\, M_{.6}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{\text{ᵛ}*}]_6\beta_6(a_6) - \mu_6(a_6))m_6(t,a_6)da_6 = ([h_*]_6\beta_6(\theta_6) - \mu_6(\theta_6))M_{,6}$$

Where $0< \ \theta_6 = \theta_6(t) < \infty$.

And if $a_6^2\,[a_6^2]_6$<4 $[a_2]_6$<4 $a_2$e,

Re $\lambda_6\left[\lambda_6\right]_6 = \text{Re}\lambda_2\left[\lambda_2\right]_6 = -\frac{6}{2}a_6[\text{a2}\ ]_6 < 0.$

Let ⊿=max$_{2<u<\infty}\left|\text{h}_*\,[\text{h}_*]_6\beta_6(a_6) - \mu_6(a_6)\right|$. Then$|n_6(\text{t})|$ ⊿∀t . Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium($\rho^*(\text{a}), \boldsymbol{Q}^*$) is asymptotically stable if either $a_6^2$ <4 $a_2$ e and $[a_6] > 2\Delta$,or $[a_6^2]$ ⊿$[a_2]$e and $(a_6 - \sqrt{(a_6^2)_6 - 4(a_2)_6 e}) > 2\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+6}\right)$ and $\beta_6(a_6) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_6(a_6) e^{-\int_0^a \mu_6(\tau_6) + d\tau_6 da} = \frac{6}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_6 + [\mu_0]$ 0.5. Assume $[f_6](x_6) = K_6 x_6$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_6(x_6) = [0x_6]/(1 + x_6)$; the functional responses $[[f_2](x) = [0]\, x_6/(1 + x_6)$ and $h_6(x_6)=(1+\omega_6 x_6)/(1 + x_6)$ (See Holling [7] and Rashevsky [13])

Equation (64),now has the form of

$$\frac{m_6 + [\mu_0] + [K][P_0])}{[m_6]_6 + (\mu_0)_6 + [K][P_0])^3_6\ [K]_6 [P_0]^3_6} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(6 + P_0) < b_6$ i e $[P_0] < b/(\delta_6 - b_6)$ then the equilibrium $(P_0)\ [P_0]_6(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_6$

$[f_{6*}] = K_6\ P^*\ [f_{2*}] = \frac{\sigma Q_6^*}{6+Q_6^*} \quad [h_*] = \frac{6+\omega_6 Q_6^*}{6+Q_6^*}$

$[f_{6*}] = K_6\ [f^6_{2*}] = \frac{\sigma_6}{(6+Q_6^*)^2} [\ g^6_*] = \frac{\sigma_6}{(6+P_6^*)^2} \quad [h^6_*] = \frac{\omega_6 - 6}{(6+Q_6^*)^2}$

In order to have (H3) and (H4) hold, $m, \mu_0$, K, b, $\delta$, and $\omega$ need to satisfy

$$\frac{\omega_6}{(m+[\mu_0]_6+1+\frac{K_6b_6}{\delta_6+b_6}+\sigma)^2}+\frac{2\omega_6}{([m_6]_6+[\mu_0]_6+1+\frac{K_6b_6}{\delta_6+b_6}+\sigma_6)^3}<1$$

And

$$\frac{1}{(m_6+[\mu_0]_6+1+\frac{K_6b_6}{\delta_6-b_6})^2}+\frac{2\omega_6}{(m_6+[\mu_0]_6+1+\frac{K_6b_6}{\delta_6-b_6})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_6+[\mu_0]_6+1+\frac{K_6b_6}{\delta_6-b_6}+(m_6+[\mu_0]_6+\frac{K_6b_6}{\delta_6-b_6}+\sigma_6)\,Q_6^*)^3$$

$$=m_6+[\mu_0]_6+3+\frac{K_6b_6}{\delta_6-b_6}(1+Q^*)(1+\omega_6Q_6^*$$

Now

$$\left[a_6^2\right]-4\left[a_2\right]e=\left(\frac{2K_6b_6}{\delta_6-b_6}\right)^2+\frac{4K_6b_6\sigma_6\,Q_6^*}{(\delta_6-b_6)(1+Q_6^*)}+\frac{\sigma_6^2Q_6^{*2}}{(1+Q_6^*)^2}$$

$$\frac{4\,b_6\sigma_6\,(\delta_6-b)Q^*}{\delta_6(1+Q_6^*)^2}-\frac{4\,(\delta_6-b_6)^2(1-\omega_6)Q_6^*}{\delta_6\left(1+\omega_6Q_6^*\right)1+Q_6^*)}P^*(0)$$

And

$\Delta = \lim_{0<a<\infty} |[h_*]_6 ae^{mu}_6 \ [\mu_0]_6 - \frac{a_6}{a_6+6}$

Hence if $[a^2_6 > 4(a_2)_6 e$ and

$$[a_6] = \frac{2K_6 b_6}{\delta_6 - b_6} + \frac{\sigma_6 Q^*_6}{6 + Q^*_6} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_6]_6 = [f_{6*}]_6 + [f_{2*}]_6 + [f_{6*}] P^*_6$$

$$[a_2]_6 = \left[f^6_{2*}\right] P^*_6 - \frac{h^6_*}{h_*} P^*(0)$$

$$e=[g^6_*]_6 Q^*_6$$

$$\eta_6(t) = \left[h_*\right]_6 \beta_6 \left(\theta_6(t)\right) - \mu_6(\theta(t)),$$

The system(36)can be written as the following system:

$$[M(t)]\left[a_6\right]M_6 - [a_2]_6 W_6 - [n(t)]M_6$$

$$w_6(t) = eM_6$$

Write

$$A=\begin{pmatrix} -[a_6]_6 - [a_2]_6 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_6 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_6 \\ W_6 \end{pmatrix}$$

The solution of (32)is

$$X_6(t) = e^{A,t}[X]_6(0) + \int_0 e^{At}$$

Other problems to which the model is applied are: To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case β→∞ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness

condition of (e) entropy. It is shown that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure

authors have discovered. It is shown that Legendre invariant thermodynamic scalar curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws, leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions, geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged

AdS black holes by considering the case of the RN-AdS black holes surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase transition. For the case $\omega_q=-2/3$ authors derive analytic expressions of (e) critical physical quantities, while (e) for cases $\omega_q\neq-2/3$ appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits ($e>0.96$) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE]

(or arXiv:1212.3342v1 [astro-ph.HE] for this version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=) chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum

Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed

quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions, while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma? Paul M. Chesler STOP Paul M. Chesler, Krishna Rajagopal** calculate how the energy and (e&eb) the opening angle of jets in N=4 SYM theory evolve as they propagate through (e&eb) the strongly coupled plasma of that theory. **Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They

obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb) heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version)

**Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb) strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar

(=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI: 10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills (SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-

equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In

this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e)

pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and**

**Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April**

**2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its

non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstates thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11**

**August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent

information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.** Characterizing the behaviour of strongly coupled quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal

fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at λ=λc, the correlation length diverges as ξ∼|λ−λc|−ν, while (e) the gap between the ground state and first excited state vanishes as (e) ΔE∼ξ−z∼|λ−λc|zν. Exactly at the critical point λ=λc, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as ω∼|k−k0|z. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004**

**STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element. The basic architecture presented here can be expanded, offering (eb) the possibility for the coherent interaction of (e&eb) a large number of superconducting qubits. **Nature 449, 438-442 (27 September 2007) | doi:10.1038/nature06124; Received 18 April 2007; Accepted 25 July 2007 Coherent quantum state storage and transfer between two phase qubits via a resonant cavity Mika A. Sillanpää1, Jae I. Park1 & Raymond W. Simmonds1 STOP. Chui-Ping Yang, Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no

transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e) long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007**

## PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION

## ATTRITION MODELS

## MOODULE NUMBER SEVEN

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are**

**assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Systems to which have applied the model are as follows: Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures have been proposed. An attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in (eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP. In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin**

**qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011 http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 Qubit transfer between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys. Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer. Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic**

**operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP. In recent years S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P have investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys. Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP. Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2 [quant-ph] for this version) Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP. Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue from over (e&eb) a thousand plant and animal samples. NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP. Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this**

**behavior, enhancing and accelerating (eb, eb+) the growth of correlations. Michael Kastner gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. Michael Kastner tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP. A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E. Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91, 075435 – Published 26 February 2015 STOP. Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows**

**(e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1 proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point'' with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that (eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS**

**metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035 Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013 STOP Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen STOP Authors study the dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and**

**Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) Dynamic Service**

**Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by H∗(X,(Λ∗TX)⊗g), where X is the target hyperk\"ahler manifold. We further give a mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition.**

**Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb) the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537 © 2016 American Physical Society Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant (Δα/α) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and Δα/α to (e&eb) this conclusion. However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts z≲3, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the Λ cold dark matter model value (jΛCDM=1); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in α(z), linked to (e&eb) a linear coupling (ζ) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order |ζ|∼10−4, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in Δα/α (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical**

**Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems, authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP Authors measure the star formation properties of (e) two large samples of galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the star formation rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a function of environment. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]. Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size √ n × √ n? The problem is that √n time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time O (√n) for d ≥ 3, or O (√nlog5/2n) for d = 2. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among**

**several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ⁎This work was mostly done while the author was a PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship and by CIAR. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e) sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space**

**regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP The relationship between the foundations of mathematics and (e&eb) physics is a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb) flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA**

**conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph] (or arXiv: 1503.05600v1 [math-ph] for this version) Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to the fine structure constant. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling**

**for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor goes to 0 as the (e&eb) time goes to 0, the big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph]**

**(or arXiv: 1306.4613v1 [quant-ph] for this version) Effects of gauge theory based number scaling on geometry Paul Benioff STOP In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x, outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph] (or arXiv: 1211.3381v2 [quant-ph] for this version) Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show**

**that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value, \int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances, the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include (e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP This work is based on (e) the field of reference frames based on (e) quantum representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical**

**systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul Benioff STOP A quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and**

**rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference: Journal of Physics: Conference Series 70 Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows: System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol**

## 14. MODEL FORMULATION

Let $p_7(t, a_7)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(g_7)Q_7$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_7(0) = 0, \quad g_7(P_7) > 0, \quad P_7 \ 0, \qquad \lim_{P_7 \to \infty} g_7(P_7) = (K_7)_7 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_7 + (\rho_a)_7 + D_7(a_7, P_7, Q_7)\rho_7 + 0, \qquad 0 < a_7 \ \infty,$$

$$\rho_7(t,0) = \int_0^{\infty} n_7(a_7, P_7, Q_7)\rho_7(t_7, a_7)da_7$$

Where $D_7$the mortality or debit type function is function and $n_7$ is the fertility function.

Put $_7$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_7}{dt}\right]_7 = x_7\left[\frac{dQ_7}{dt}\right]_7 - x_7\left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_7}{dt}\right]_7 = x_3\left[\frac{dQ_3}{dt}\right]_7 - x_4\left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_7}{dt}\right]_7 - x_6\left[\frac{dQ_3}{dt}\right]_7$$

Where $x_7, x_7, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_7$ and $g_7(P_7)Q_7$

Now when you add the term $(b_7Q_7)$ with $g_7(P_7)(Q_7)$ we are infect see is that $(b_7Q_7)$ is credited incrementally part of $g_7(P_7)(Q_7)$ or instant nearly, when subtracted from $\frac{dQ_7}{dt}$ , it means that $g_7(P_7)(Q_7)$ is slowly eaten by $\frac{dQ_7}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

Following equation throws light on important issues in the following cases

**Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA**

(946)(Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor.

(947)In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface.

(948)As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone.

(949) Such planets are all the "right" distance from (e) their star, and hopefully, also have water.

(950)But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone.

(951)For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement.

(952) But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway.

(953) But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain.

(954)It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life.

(955)Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface.

(956)The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113

**Abstract**

(957)The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question.

(958)Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry.

(959)This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry.

(960)Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments.

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| (961)Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy. |
| (962)The electronic and (e&eb) structural variability among them is extraordinary. |
| (963)The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations. |
| (964)The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface. |
| (965) The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers. |
| (966)Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp |

**Water waves travel back in time to retrace their ripples By Jacob Aron (New Scientist)**

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| (967)Nudge your bath tub just right and (eb) you might end up reversing time. |
| (968)Don't worry – you won't end up wet, naked and without a towel in the Cretaceous period. Emmanuel Fort at the Langevin Institute in Paris, France, and his colleagues have developed what they call a "time mirror", which can make water waves rewind into (e&eb) their past. Humans and rubber duckies are (=) affected. |
| (969)To understand how this works, imagine dropping a stone in a pond and watching waves ripple out from a central point in an ever-expanding circle. The equations governing the behaviour of waves say that it should also be possible for a ring of ripples to refocus back into a single point, like a reversed movie of the dropped stone. But we never actually see this happening, because it is vanishingly unlikely that an outside-in ring would form spontaneously. Put it in reverse |
| (970)Fort and others have previously shown that it is possible to engineer the situation so the waves turn back on themselves by kitting out the edge of your pond with sensors and wave generators, effectively recording and then playing back the waves in reverse. The process has been done with water, sound and even microwaves, and is used to produce clearer ultrasound images for medical uses. |
| (971)But now the team has discovered that simply dropping the tub downwards is enough to send water waves scurrying in reverse, with no need for complex sensors. That's because the velocity of water waves depends on acceleration due to gravity. "If you do a sudden jolt, you change the effective gravity, and the velocity of the wave changes," says Fort. |
| (972)The disruption essentially resets the starting point of the wave. Once normal gravity is restored, it starts to travel both backwards and forwards. "Suddenly, you freeze the wave. When you release it again, it doesn't remember that it was going one way," says Fort. |
| (973)Ripples in the water returned exactly to their start point - the shape of the Eiffel Tower |
| (974)Ripples in the water returned exactly to their start point – the shape of the Eiffel Tower |

**Déjà vu**

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| (975)To show that the backward-travelling waves converge exactly on their starting point, just as if you were rewinding time, the team conducted a series of experiments. In some, they simply created a circular ripple of waves by dipping a probe in the water, similar to dropping a stone in the pond. In others, they used wave emitters in the shape of France, the Eiffel Tower and a smiley face to create complex wave-fronts. In each case, the reversed waves converged to recreate their original shape. |
| (976)The experiments aren't truly reversing time, because the waves continue to travel out again once they have converged at the central point, points out Fort. "A real time-reversal would have a sink that absorbs all the waves," he says. |
| (977)Now the team hopes to conduct similar experiments with other kinds of waves, like light or sound, although it may prove more difficult to find an analogue of their time-mirror jolt in other wave mediums, he says. |
| (978)"The work is spectacular," says Geoffrey Edelmann of the US Naval Research Laboratory in Washington, D.C. "They essentially change the rules that govern how the wave propagates and, in doing so, create an entirely new way to focus waves free of spatial constraints." He also thinks the technique could be extended to other kinds of waves, and perhaps have medical applications such as lithotripsy, which uses ultrasound waves to break up kidney stones. Journal reference: Nature Physics, DOI: 10.1038/nphys3810 |

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| **Time reversal and holography with spacetime transformations Vincent Bacot, Matthieu Labousse, Antonin Eddi, Mathias Fink & Emmanuel Fort Nature Physics (2016) doi: 10.1038/nphys3810** |

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| (979)Wave control is usually performed by (e) spatially engineering the properties of a medium. |
| (980) Because time and space play similar roles (e&eb) in wave propagation, manipulating time boundaries provides (eb) a complementary approach. |
| (981) Here, authors experimentally demonstrate the relevance of this concept by (e) introducing instantaneous time mirrors. |
| (982)Authors show with water waves that a sudden change of (e) the effective gravity generates (eb) time-reversed waves that refocus at the source. |
| (983)They generalize this concept for (e) all kinds of waves, introducing (eb) a universal framework which explains (eb) the effect of any time disruption on (e&eb) wave propagation. |
| (984)Authors show that sudden changes of the medium properties generate (eb) instant wave sources that emerge (eb) instantaneously from the entire space at (eb) the time disruption. |
| (985)The time-reversed waves originate from (e) these 'Cauchy sources', which are (+) the counterpart of Huygens virtual sources on (eb) a time boundary. It allows us to revisit the holographic method and introduce a new approach for wave control. |

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| **Viewpoint: A Bird's Eye View of Circuit Photons Sebastian Schmidt, Institute for Theoretical Physics, ETH Zurich, 8093 Zurich, Switzerland June 30, 2016• Physics 9, 74 A scanning probe detects the quantum states of photons in a microwave circuit, providing the information needed for quantum simulations. Figure caption expand figure A. Houck/Princeton (deleted)** |

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| (986)Figure 1: Scanning defect microscopy provides a map of photons in a resonator lattice. Houck and colleagues demonstrated the technique using (e) 49 resonators (grey lines) that were coupled together to (e&eb) form a kagome lattice. |
| (987)Quantum mechanics rules the dynamics of (e) light and matter. |
| (988) Yet performing a quantum-mechanical simulation of a material from first principles is (=) practically impossible on a classical computer because (e) the complexity of the simulation increases (eb) exponentially with the number of particles involved. |
| (989) The solution, according to Richard Feynman, was to build a machine out of (e) quantum building blocks that could directly emulate (e&eb) the material itself [1]. |
| (990)Prototypes of such quantum simulators that are based on (e) ultracold atoms, ions, photons, and superconducting microwave circuits are now available [2], with (e&eb) the latter, in particular, having attracted (e) Silicon Valley's interest. |
| (991)The challenge with these circuit-based simulators, however, is that they are (=) 2D, which complicates the readout of their constituent elements. Andrew Houck from Princeton University, New Jersey, and colleagues have now delivered an attractive solution by developing a technique [3], called scanning defect microscopy (Fig. 1), that determines (eb) the number of photons occupying (e) each mode of a 2D microwave circuit. |
| (992)It is this information that would serve as (=) the fundamental input and output for (e) certain quantum simulations. |
| (993)Superconducting microwave circuits combine electronic and (e&eb) photonic degrees of freedom [4, 5]. |
| (994)The main element of the circuit is (=) a transmission line, which is made up of (e) a central superconducting wire separated by a gap from (e) two grounded plates. |
| (995)All of these structures are on a single plane, as if one had taken a 2D slice through (e&eb) a coaxial cable. |
| (996)When truncated, the transmission line becomes (=) a resonator, which can host (e&eb) discrete photon modes within (eb) its gaps |
| (997) Large lattices of resonators can be engineered in various 1D or 2D geometries by coupling (e&eb) two, three, or more resonators together via (e&eb) a capacitive interface. In many ways, photons in such devices behave similarly to electrons in a solid. |
| (998)To make microwave circuits that can simulate (e&eb) quantum phenomena faster than a classical computer, however, resonator lattices have to be integrated with (e&eb) superconducting qubits |
| (999). Such qubits are controlled with (e&eb) electrical currents in a Josephson tunnel junction, and in many respects, they behave (e&eb) like artificial atoms, which couple to (e&eb) the photons in the circuit. |
| (1000) As a result, superconducting microwave circuits can be used to (e) explore the coupling between the quantum states of light and (e&eb) matter, the regime of (e) circuit quantum electrodynamics (cQED). |
| (1001) Photons in these devices often exhibit (eb) striking matter-like behavior [6, 7], providing (eb) the basis for the simulation of complex materials. |
| (1002) Such circuits can be fabricated on a substrate using (e) standard lithographic techniques, with (e&eb) qubits and resonators that are (=) hundreds of micrometers or even millimeters in size. |
| (1003) Small-scale realizations of cQED lattices have already been used for (e) proof-of-principle quantum simulations [8, 9]. |
| (1004) However, future large-scale simulations will have to be initialized and measured (eb) efficiently—steps known as (=) drive and readout. |

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| (1005) This requires connecting the lattices to (e) external circuitry, such as voltage sources, circulators, and amplifiers, which typically reside outside (e) the chip. |
| (1006) So far, this connectivity has been realized through (e&eb) external transmission lines, but these elements can only couple to (e&eb) the resonators at the edges of (e) the lattice because of (e) its 2D nature. |
| (1007) Such a side view provides (eb) only limited information about (e&eb) the quantum state and spatial correlations of photons inside (eb) the bulk resonators, seriously limiting (e) the applicability of circuit QED devices for (e) quantum simulations. |
| (1008) The probe developed by Houck and colleagues, scanning defect microscopy (SDM), solves (e) this problem by (e) moving outside the plane of the chip to provide (eb) a fuller image of photons within (eb) the chip. |
| (1009) In SDM, a small sapphire tip is positioned a few tens of (e) micrometers above the surface of one of the resonators (Fig. 1). |
| (1010) The presence of this dielectric modifies (e&eb) the evanescent electric field from the circuit photons, which extends (eb) just outside the plane of the chip. |
| (1011) Specifically, the tip increases (eb) the local capacitance of the transmission line, downshifting (e) the local resonator frequency. |
| (1012) This in turn causes (eb) a drop in the transmission of photons through (e&eb) the lattice when (e) it is driven at its nominal resonance frequency. |
| (1013) Houck's team showed that if the sapphire tip has only a weak (linear) effect on the resonator frequency, then (eb) a simple analytic relationship exists between the induced frequency shift, (e&eb) the change in transmission, and (e&eb) the local photonic occupation of each resonator. |
| (1014) This relationship provides (eb) a recipe for detecting the local photon occupancy inside (eb) the lattice. |
| (1015) First, calibrate the probe by determining (eb) the frequency shift for (e) a single isolated resonator as (=) a function of the tip's position. |
| (1016) Second, measure the transmission through (e&eb) the whole lattice with and without (e&eb) the probe. |
| (1017) Finally, use the measured transmission drop with (e&eb) the probe present to find (eb) the photonic occupation of (e) an individual resonator in (eb) the lattice. |
| (1018) A map of the photon-mode occupancies can be obtained by (e) scanning the probe across the whole lattice. |
| (1019) Houck's team used (e) SDM to image the occupancies of (e) the photon modes of (e) 49 coupled resonators arranged in (eb) a special geometry called a kagome lattice (see details in Fig. 1). |
| (1020) (To simplify the measurement, no superconducting qubits were (=) incorporated with (e&eb) the resonators.) |
| (1021) The occupancies they measured were in excellent agreement with (=) theoretical predictions, corresponding to (e&eb) a fidelity of 0.99. |
| (1022) The high fidelity is even more impressive considering the weak signal-to-noise ratio of (e) their technique (which is limited by (e) the intentionally small size of the frequency-shifted region) and (e&eb) the technical challenge of (e) positioning the probe inside (eb) a cryostat for (e) in situ measurements. |
| (1023) While SDM shares similarities with (e&eb) so-called bead-pull experiments for accelerators and scanning gate microscopy for (e) nanostructures, it is truly a (=) new tool for quantum electronics. |

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| (1024) First application of SDM is likely to be (=) the simulation of 2D topological phenomena and (e&eb) geometrically frustrated lattices (of which the kagome lattice is (=) one example.) |
| (1025) Both are associated with (e&eb) macroscopically degenerate systems, which are (=) notoriously difficult to simulate on (e&eb) a classical computer. |
| (1026) Such circuits are (=) genuinely open systems (photons can be added or lost). |
| (1027) This combined with (e&eb) their inherent strong light-matter coupling, makes (eb) such quantum simulators ideal candidates for (e) the study of many-body, out-of-equilibrium phenomena, such as (=) thermalization or nonequilibrium phase transitions. |
| (1028) SDM could also lead to (eb) the discovery of new exotic phenomena, which are (=) unknown in traditional solid-state systems. |
| (1029) In a nutshell, SDM constitutes a new and indispensible tool for (e) future experiments on quantum simulations with microwave photons. This research is published in Physical Review X. |

**Imaging Photon Lattice States by Scanning Defect Microscopy D. L. Underwood, W. E. Shanks, Andy C. Y. Li, Lamia Ateshian, Jens Koch, and A. A. Houck Phys. Rev. X 6, 021044 – Published 30 June 2016**

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| (1030) Microwave photons inside (eb) lattices of coupled resonators and superconducting qubits can exhibit (eb) surprising matter like behavior. |
| (1031) Realizing such open-system quantum simulators presents (eb) an experimental challenge and requires (e) new tools and measurement techniques. |
| (1032) Here, authors introduce scanning defect microscopy as one such tool and illustrate its use in mapping (e&eb) the normal-mode structure of microwave photons inside (eb) a 49-site kagome lattice of coplanar waveguide resonators |
| (1033) Scanning is accomplished by (e) moving a probe equipped with a sapphire tip across (e&eb) the lattice. |
| (1034) This locally perturbs (e&eb) resonator frequencies and induces (eb) shifts of the lattice resonance frequencies, which authors determine by (e) measuring the transmission spectrum. |
| (1035) From the magnitude of mode shifts, author can reconstruct (eb) photon field amplitudes at (eb) each lattice site and thus create (eb) spatial images of the photon-lattice normal modes. Received 4 December 2015 DOI:http://dx.doi.org/10.1103/PhysRevX.6.021044 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society |

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| We assume that |
| $$D_7 = \mu_7(a_7) + (f_7)_7(P_7)_7 + (f_7)_7(Q)_7 \quad (7)$$ |
| $$n_7 = \beta_7(a_7)h_7(Q_7), \quad (3)$$ |

Where $\mu_7$the age- dependent debit rate whit is $\mu_7(a_7)$ $\forall a > 0$; $(f_7)_7$ is the density- dependent death rate with $f_7(0) = 0$, $(f_7^7)_7(P_7)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_7)_7(P_7) = (k_7)_7$ ; $f_7$ is the functional response of creditors to prey with $[f_7]_7(0) = 0, (f_7^7)_7(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_7)_7(Q_7) = [k_7]$ ; $\beta$ is the age dependent birth rate with $\beta(a)$ $0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^7(Q)$ $0, \forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_7 + (\rho_a)_7 = -\mu_7(a_7) + (f_7)_7 P_7) + \big(f_7)_7(Q_7)\big)(\rho_7)_7 \quad (4)$$

$$\rho_7(t_7, 0) = h_7(Q_7)\int_0^\infty \beta_7(a_7)\rho_7(t, a_7)da, \quad (5)$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(P_7)Q_7 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_7$ and $Q_7$ of (4-6) satisfy

$$0 \ \rho_7(t, a_7) \ \rho_7^*(t, a_7),$$

$$0 \ Q_7(t) \ Q_7^*(t),$$

Where $\rho_7^*(t, a_7)$ and $Q_7^*(t)$, are solutions of

$$(\rho_t^*)_7 + (\rho_a^*)_7 = -(\mu_7(a_7) + (f_7)_7(P_7^*) + (f_7)_7 Q_7^*))\,\rho_7^*, \quad (7)$$

$$\rho_7^*(t, 0) = \int_0^\infty \beta_7(a_7)\,\rho_7^*(t, a_7)da_7 \quad (8)$$

$$\frac{dQ_7^*}{dt} = -b_7\, Q_7^* + g_7(P_7^*)Q_7^*, \quad (9)$$

With $P_7^*(t) = \int_0^\infty \rho_7^*\ (t, a_7)da_7$

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From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_7^*(t)}{dt} = P_7^*(t)\big(r_7 - (f_7)(P_7^*) - (f_7)_7(Q^*)\big), \quad (10)$$

$$\frac{dQ_7^*(t)}{dt} = Q_7^*(t)(-b_7 + g_7(P_7^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_7(a)e^{-\int_0^a \mu_7(\tau)dt}e^{-ra}\, da_7 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_7^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(t)dt}\, da_7 > 1, \quad (H1)$$

Which guarantees $r_7 > 0$.

*Remark.* If $\beta_7(a) = [\beta_0]_7 a_7 e^{-\alpha_7 a_7}, \mu_7(a_7) \equiv \mu_7$ then the solution of (12) is

$$r_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_7]_7(P_7) = 0, [f_7](Q_7) = \lambda_7 Q_7 \; h_7(Q_7) = 1, \qquad and \;\; g_7(P_7) = c_7 P_7$$

$$[\rho_7]_7 + [\rho_a]_7 = -(\mu_7 + \lambda_7 Q_7)\rho_7$$

$$\rho_7(t,0) = \int_0^\infty [\beta_0]_7 a_7 e^{-\alpha_7 a_7} \rho_7(t, a_7)\, da_7$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + c_7 P_7 Q_7$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_7}{dt} = \omega_7 P_7 - \lambda_7 P_7 Q_7$$

$$\frac{dQ_7}{dt} = b_7 Q_7 - c_7 P_7 Q_7$$

Where $\omega_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$

## 15. THE EQUILIBRIUM $(\rho_0)_7(a_7), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_7(a_7), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_7 + (\rho_a)_7 = -\big(\mu_7(a_7) + (f_7)_7(P_7)\big)\rho,$$

$$\rho_7(t,0) = \int_0^\infty \beta_7(a_7)\rho_7(t,a_7)da_7 \qquad (13)$$

The solution of (13) has the form of

$$\rho_7(t,a_7) = Ce_7^{r(t-a_7)} - \int_0^a (\tau)d\tau - \int_0^7 (f_7)_7\big(P_7(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_7(t) = C_7 e^{r_7 t - \int_0^t (f_7)_7(P_7(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_7^{-7}(r)$ if $k_7 > r$. This is equivalent to

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{(k_7)_7 a_7} da_7 < 1.$$

In fact, if we define

$$R_7(s_7) = \int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau)d\tau} e^{-s_7 a_7}\, da_7$$

$R(s)$ is a decreasing function. $R_7(r_7) = 1$ And $R_7(k_7)_7 < 1$ imply $r_7 < (k_7)_7$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_7(a_7), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{k_7 a_7} da_7$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_7 \equiv \int_0^\infty (\rho_0)_{77}(a_7)da_7$ *satisfies*

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{-(f_7)_7(P_0)a_7} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_7(a_7), 0)$

Let $u_7(t, a_7) = \rho(t, a_7) - \rho_0(a_7)$ Then

$$N_7(t) \equiv \int_0^\infty u_7(t, a_7)da_7 = \int_0^\infty \rho_7(t, a_7)da_7 - \int_0^\infty (\rho_0)_7(a_7)da = P_7(t) - P_0)_7$$

And

$$(u_7)_7 + (u_a)_7 = (\rho_t)_7 + (\rho_a)_7 - (\rho_0^7)_7(a_7)$$

$$= -\mu_7(a_7) + (f_7)_7(P_7) + (f_7)_7(Q_7))(u_7 + (\rho_0)_7) + \mu_7(a_7)(\rho_0)_7 + (f_7)_7(P_0)_7(\rho_0)_7$$

$$= -\mu_7(a_7)u_7 - (f_7)_7(P_7) + (f_7)_7(Q_7))u_7 - (\rho_0)_7(f_7)_7(P_7) - (f_7)_7(P_0)_7$$

$$+ (f_7)_7(Q_7)), \qquad (15)$$

Where $(\rho_0^7)_7(a_7)$ denotes $d(\rho_0)_7(a_7)/da_7$

Expansion $(f_7)_7(P_7)$ and $(f_7)_7(Q_7)$ at $(P_0)_7$ and 0 yield respectively

$$(f_7)_7(P_7) = (f_7)_7((P_0)_7 + N_7)$$

$$= (f_7)_7(P_0)_7 + (f_7^7)_7(P_0)_7 N_7 + o_7(N_7), \qquad (16)$$

$$(f_7)_7(Q_7)$$

$$= (f_7^7)_7(0)Q + o_7(Q_7) \qquad (17)$$

Write $f_7(P_0) = f_{70}, f_7'(P_0) = f_{70}'$, and $f_7'(0) = f_{70}'$ and substitute (16) and (17) into (15) to obtain

$$(u_7)_7 + (u_a)_7 \approx -(\mu_7(a_7) + (f_{70})_7)u - (f_{70}^7)_7(\rho_0)_7(N_7) - (f_{70}^7)_7(\rho_0)_7(Q_7)$$

$$u_7(t,0) = \rho_7(t,0)(\rho_0)_7(0)$$

$$= h_7(Q_7)\int_0^\infty \beta_7(a_7)(u_7(t,a) + (\rho_0)_7(a_7)da_7 - (\rho_0)_7(0)$$

$$= h_7(Q_7)\int_0^\infty \beta_7(a_7)(u_7(t,a_7)da_7 + (h_7(Q_7) - 1)(\rho_0)_7(0) \quad (18)$$

Expanding

$$h_7(Q_7) = 1 + (h^7)_7(0)Q_7 + 0(Q_7) \equiv 1 + h_0'Q_7 + 0Q_7$$

And substituting it into (18), we arrive at

$$u_7(t,0) \approx \int_0^\infty \beta_7(a_7)u_7(t,a_7)da_7 + (h_{a_7}')(\rho_0)_7Q_7$$

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PREDATOR-PREY POPULATION MODELS

The linearization is the following :

$$[u_7] + [u_7] = -(\mu_7(a_7) + [t_{70}]_7u_7 - [\alpha_7]_7(a_7)N_7 - [\alpha_7]_7(a_7)Q_7$$

$$u_7(t,0) = \int_0^\infty \beta_7(a_7)u_7(t,a_7)da_7 + [\xi_7]Q_7 \quad (19)$$

$$\frac{dQ_7}{dt} = -b_7Q_7$$

where

$$[\alpha_7]_7(a_7) = [t_7^7 a]_7[p, o]_7(a_7)d2(a_7) = [f_{70}^7]_7(\rho\, o\,)_7(a_7)\xi = [\mathrm{h\ o}^7]_7[\rho \mathrm{o}]_7(\mathrm{o}) \text{ and } \mathrm{b}_7 = \mathrm{b}_7 - g_7(P_7)$$

If $g_7(P_0) > b$, clearly, $(\rho_0)_7(a_7), 0)$ is unstable and the predator population is persistent. Assume $g_7(P_0)_7 < b$.

Then, $Q_7(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_7 + u_a = -(\mu_7(a_7) + (f_{70})_7)u_7 - (\alpha_7)_7(a_7)N_7$$

$$= -(\mu_7(a_7) + f_{70})u_7 - \alpha_7(a_7)\int_0^\infty u_7(t, a_7)da_7$$

$$u_7(t, 0) = \int_0^\infty \beta_7\,(a_7)u_7(t, a_7)da_7$$

Using the method developed by Cushing [2, 3], we suppose

$$u_7(t, a_7) = \tilde{\mathrm{u}}_7(a_7)^{e^{\gamma(t-a_7)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_7}\, da_7$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_7)$ gives,

$$\tilde{\mathrm{u}}_7(a_7) = e^{-\int_0^a \mu_7(\tau_7)d\tau - (f_{70})_7 a}\left(\tilde{\mathrm{u}}_7(0) - w\int_0^a e^{\int_0^\tau \mu_7(s_7)ds + (f_{70})}\tau_7\,(\alpha_7)_7(\tau_7)e^{\gamma_7\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_7(a_7)e^{-\gamma_7 a_7 - (f_{70})_7 a_7 - \int_0^a \mu_7(\tau_7)d\tau}\right)\tilde{\mathrm{u}}_7(0)$$

$$+w\int_0^{\infty}\beta_7(a_7)e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau}\int_0^a(\alpha_7)_7(\tau_7)e^{-\gamma_7\tau_7}e^{(f_{70})_7\tau_7+\int_0^{\tau}\mu_7(s_7)ds}d\tau_7\,da_7$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^{\infty}e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau}da\right)(\tilde{u}_7(0))_7$$

$$+\left(-1-\int_0^{\infty}e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau}\int_0^a e^{(f_{70})_7\tau_7+\int_0^{\tau}\mu_7(s_7)ds+y\tau}\alpha_7(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(t_7)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_7\,(a_7)Y_7(\gamma_7,a_7)da_7\right)u_7(0)$$

$$+w\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\,da_7=0,$$

(24)

$$\left(\int_0^{\infty}Y_7(\gamma_7,a_7)da_7\right)\tilde{u}_7(0)$$

$$+\left(-1-\int_0^{\infty}Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\,da_7\right)w=0.$$

Equation (20) has a solution of the form of $u_7(t,a_7) = u_7(a_7)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_7(0) \not\equiv 0$ and $\gamma, w$ in the complwx plane.

Let $I_7(\gamma_7,a_7) = \int_0^a (\rho_0)_7(\tau_7)Y^{-7}(\gamma_7,a_7)d\tau_7$. We obtain the characteristic equation

$$f_{70}^{l}{}_7\left(\int_0^\infty Y_7(\gamma_7,a_7)da_7\right)\left(\int_0^\infty \beta_7(a_7)Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$+\left(1-\int_0^\infty \beta_7(a_7)Y_7(\gamma_7,a_7)da\right)\left(1+(f_{70}^{l})_7\int_0^\infty Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$= 0. \qquad (25)$$

Since

$$(\rho_0)_7(a_7) = (\rho_0)_7(0)e^{-\int_0^a \mu_7(\tau_7)d\tau_7-(f_{70})_7a_7}$$

$$I_7(\gamma_7,a_7) = \int_0^a (\rho_0)_7(\tau_7)e^{\gamma_7\tau_7+(f_{70})_7\tau+\int_0^\tau \mu_7(s_7)ds_7}d\tau$$

$$= \frac{(\rho_0)_7(0)}{\gamma_7}(e^{\gamma_7a_7}-1)$$

Let

$$E_7(a_7) = e^{-(f_{70})_7a_7-\int_0^a \mu_7(\tau_7)d\tau_7}$$

Then $[Y_7]_7[\gamma_7]_7,a_7) = E_7(a_7)e^{-\gamma_7a_7}$. by a simple calculation,

$$\int_0^\infty \beta_7(a_7)Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da = \frac{[\rho_0]_7(0)}{\gamma}\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}\ da\right);$$

$$\int_0^\infty Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da = \frac{[\rho_0]_7(0)}{\gamma_7}\left(\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}\ da_7\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_7(0)}{\gamma_7}[f_{70}^7]\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\ (\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7)$$

$$+\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right).$$

$$\left(1+\frac{(\rho_0)_7(0)}{\gamma_7}(f_{70}^{'})_7\left(\int_0^\infty E_7(a_7)da_7-\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right)(1+\frac{(f_{70}^{'})_7}{\gamma_7}(\rho_0)_7(0)\int_0^\infty E_7(a_7)da_7)=0.$$

We need ũ(a) $\not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)da_7 = 1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7 \neq 0,$$

Since $\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)\ \int_0^a \alpha_7(\tau_7)Y^{-7}(\gamma_7,\tau_7)d\tau_7 da_7 > 0$. Therefore, the characteristic solution is

$$\gamma_7 = -(f_{70}^{'})_7(\rho_0)_7 \int_0^{\infty} E_7(a_7)da_7.$$

Since $f_7'(P_0)_7 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_7), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE OF DEBIT AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_7), 0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_7 + (\mu_7(a_7) + [f_7]_7(P_7^*) + [f_7]_7(Q_7^*))P_7^* = 0, \quad (76)$$

$$p_7^*(0) = h_7(Q_7^*)\int_0^{\infty} \beta_7(a_7)p_7^*(a_j)p_7^*(a_7)da_7 \quad (77)$$

$$g_7(P_7^*) - b_7 = 0, \quad (78)$$

Where $P_7^* = \int_0^{\infty} p_7^*(a_7)\, da_7$

Equation (28) yields $P^* = g^{-7}(b)$,

From (26),

$$p_7^*(a) = p_7^*(0)e^{\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(Q_7^*)a_7}$$

Substituting into (27), we need to solve

$$\left(1 - h_7(Q_7^*)\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(Q_7^*)a_7}da_7\right)p_7^*(0) = 0$$

For $Q_7^*$

Fix $p_7^*(0)$ and $P_7^*$ Let

$$H_7(z_7) = 1 - h_7(z_7)\int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(z_7)a_7}da_{,7}$$

And assume that

$$[k_3]_7\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(g_7^{-7}(b_7))a_7 - [k_7]_7 a_7}\, da_7 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_7(z_7)] > 0.$$

Clearly,

$$\frac{dH_7}{dz_7} = h_7(z)(\,[f_7^{'}]_7(z_7)\,)\int_0^\infty a_7\beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - f_7(P_7^*)a_7 - f_7(z_7)a_7}da_{7}$$

$-h'(z_7)\int_0^{\infty}\beta_7(a_7)\,e^{\int_0^a \mu_7(\tau_7)d\tau_7-[f_7]_7(P_7^*)a_7-[f_7]_7(z_7)a_7}\quad da_7$

(See author's It is a debit-credit world)

By the assumptions on $\left[f_7^7\right]_7 and\ [h^7]^7 dH_7/dz > 0$. Hence, if

$$H_7(0) = 1 - \int_0^{\infty}\beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7-[f_7]_7(P_7^*)a_7}da_7 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_7(0) < 0$ follows if $g_7^-(b_7) < [f_7^{-7}]_7(r_7)$, where $r_7$ is the solution of (12).

Indeed, since

$$R_7(s_7) = \int_0^{\infty}\beta_7(a_7)\,e^{-\int_0^a \mu_7(\tau_7)d\tau_7 e^{-y\frac{a}{7}}}da_7$$

Is a decreasing function of $s_7$ when $P_7^* = [g^{-7}]_7(b_7) < [f_7^{-7}]_7(r_7)$, that is $[f_7]_7(P_7^*) < r$, $R\left(f_7(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

**THEOREM 3**

Assume that the hypothesis (H3) holds and $R(f_7(g^{-7}(b))) > 1, i.e.,$

$$\int_0^{\infty}\beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7-[f_7]_7\left(g_7^{-7}(b)\right)a_7} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_7^*(a_7), Q_7^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_7^*(a_7), Q_7^*)$ is more complicated. We again consider perturbations $m_7(t_7 a_7) = p_7(t_7, a_7) - p_7^*(a_7)$ and $W_7(t) = Q_7(t) - Q^*$.Then

$$M_7(t) \equiv \int_0^\infty m_7(t, a_7) da_7 = P_7(t) - P_7^*,$$

$$[m_7]_7 + [m_a]_7 = [p_7]_7 + [p_a]_7 - [p^*]_7(a_7) \approx -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7$$

$$-[f_{7*}^7]_7 p_7^*(a_7)M - [f_{7*}^7]_7 p_7^*(\mathrm{a})W_7,$$

$$m_7(t,0)p_7(t,0) - p_7^*(0) \approx h_* \int_0^\infty \beta_7(a_7) m_7(t, a_7)\, da + \frac{[h_*^7]_7}{[h_*]_7}\, p_7^*(0) W_7$$

$$\frac{dW_7}{dt} = \frac{dQ_7}{dt} \approx [g_*^7]_7 Q_7^* M_{,7}$$

Where $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}^{'}]_7$ denoting $[f^{'}{}_7(P^*)]_7 [f_{7*}^{'}]_7$ denoting $[f_7]_7$ $(Q_7^*,$ $h_{7}^{*}$ denoting $h_*^{'}$ denoting $h_7^{'}$ $(Q_7^*)$, and $[g_*^{'}]_7$

Denoting $[g^{'}]_7(P_7^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_7]_7 + [m_a]_7 = -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7 - [f^{'}{}_{7*}]_7 [p_7^*](a_7)M_7 - [f_{7*}^{'}]_7 p_7^*(a_7)W_7$$

$$m_7(t,0) = [h_*]_7 \int_0^\infty \beta_7(a_7) m_7(t, a_7) da_7 + \frac{[h_*^7]_7}{[h^*]_7} [p^*]_7(0)\, W_{,7} \qquad (30)$$

$$\frac{dW_7}{dt} = [g_*^7]_7 [Q_7^*] M_7$$

Integrating $[m_7]_7(t, a_7)$ with respect to a from 0 to ∞, we have

$$\frac{dM_7(t)}{dt} = \int_0^\infty [mt]_7(t, a_7)da_7$$

$$= -\int_0^\infty [m_a]_7\,(t, a_7)da_7 - \int_0^\infty \mu(a_7)m(t, a_7)da_7$$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f'_{7*}]_7 \int_0^\infty p_7^*(a_7)da_7\right)M_7 - [f_{7*}^7]_7 \int_0^\infty p_7^*(a_7)da_7 W_7$$

$= m_7(t, 0) - \int_0^\infty \mu_7(a_7)\, m_7(t, a_7)da_7$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 - [f_{7*}^{'}]_7 P_7^* W_{.7}$$

From (30), we arrive at the equations

$$\frac{dW_7(t)}{dt} = \int_0^\infty ([h_*]_7\beta_7(a_7) - \mu_7(a_7))m_7(t, a_7)da_7$$

$$-\left([f_{7*}] + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 + \left(\frac{[h_*^{'}]_7}{[h_*]^7}\, P_7^*(0) - f_{7*}^{'} P_7^*\right)W_7,$$

(31)

$$\frac{dW_7(t)}{dt}[g'_*]_7\, Q_7^*\, M_{.7}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{n*}]_7\beta_7(a_7) - \mu_7(a_7))m_7(t,a_7)da_7 = ([h_*]_7\beta_7(\theta_7) - \mu_7(\theta_7))M_{,7}$$

Where $0< \theta_7 = \theta_7(t) < \infty$.

And if $a_7^7[a_7^7]2<4\,[a_7]2<4\,a_7e$,

Re $\lambda_7[\lambda_7]_7 = \mathrm{Re}\lambda_7[\lambda_7]_7 = -\frac{7}{7}a_7[a2\,]_7 < 0.$

Let $\Delta=\max_{7<u<\infty}\left|h_*[h_*]_7\beta_7(a_7) - \mu_7(a_7)\right|$. Then$|n_7(t)|$ ⊠⊿∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium($\rho^*(a), \boldsymbol{Q}^*$) is asymptotically stable if either $a_7^7 <4\,a_7e$ and $[a_7] > 1\Delta$,or $[a_7^7]$ ⊠[$a_7$]e and $(a_7 - \sqrt{(a_7^7)_7 - 4(a_7)_7e) > 1\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+7}\right)$ and $\beta_7(a_7) = ae^{-ma}$ Then $\Delta=[\mu_0] + 2$ Simple calculation leads to

$$\int_0^{\infty} \beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7) + d\tau_7}da = \frac{1}{(m+\mu_0+1)^7} + \frac{1}{(m+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_7 + [\mu_0]$ ⊠0.5. Assume$[f_7](x_7)= K_7x_7$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_7(x_7)= [0x_7]/(2 + x_7)$; the functional responses $[[f_7](x) = [0]\,x_7/(1 + x_7)$and $h_7(x_7)=(1+\omega_7x_7)/(1 + x_7)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_7 + [\mu_0] + [\mathrm{K}][P_0])}{[m_7]_{\,7} + (\mu_0)_{\,7} + [\mathrm{K}][P_0])_{\,7}^{\,3}\ [\mathrm{K}]_{\,7}[P_0]_{\,7}^{\,3}} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1+P_0) < b_7$ i e $[P_0] < b/(\delta_7 - b_7)$then the equilibrium$(P_0)$ $[P_0]_{\,7}(\mathrm{a}), 0)$ is stable This implies that the

cedit confederation goes to extinction but that the cedit confederation persists $b_7$

$[f_{7*}] = K_7\ \mathrm{P}^{\,*}\ [f_{7*}] = \frac{\sigma \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*} \quad [h_*] = \frac{7+\omega_7 \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*}$

$[f_{7*}] = K_7\ \left[\mathrm{f}_{\,7*}^{\,7}\right] = \frac{\sigma_7}{(7+\mathrm{Q}_7^*)^7}[\ \mathrm{g}_{\,*}^{\,7}] = \frac{\sigma_7}{(7+\mathrm{P}_7^*)^7} \quad \left[\mathrm{h}_{\,*}^{\,7}\right] = \frac{\omega_7 - 7}{(7+\mathrm{Q}_7^*)^7}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_7}{(m + [\mu_0]_{\,7} + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma)^{\,7}} + \frac{1\omega_7}{([m_7]_{\,7} + [\mu_0]_{\,7} + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma_7)^{\,3}} < 1$$

And

$$\frac{1}{(m_7 + [\mu_0]_{\,7} + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^{\,7}} + \frac{1\omega_7}{(m_7 + [\mu_0]_{\,7} + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^{\,3}} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $\mathrm{Q}^{\,*}$ which can be solved from the

following equation:

$(m_7 + [\mu_0]_{\,7} + 1 + \frac{K_7 b_7}{\delta_7 - b_7} + (m_7 + [\mu_0]_{\,7} + \frac{K_7 b_7}{\delta_7 - b_7} + \sigma_7)\ \mathrm{Q}_7^{\,*})^{\,3}$

$=m_7 + [\mu_0]_7 + 3 + \frac{K_7 b_7}{\delta_7 - b_7}(1+ \mathrm{Q}^*)(1 + \omega_7 \mathrm{Q}_7^*$

Now

$$\left[\mathrm{a}\begin{matrix}1\\1\end{matrix}\right] - 4\left[\mathrm{a}_{\,7}\right]\mathrm{e} = \left(\frac{1K_7 b_7}{\delta_7 - b_7}\right)^7 + \frac{4K_7 b_7 \sigma_7 \mathrm{Q}_7^*}{(\delta_7 - b_7)(7 + \mathrm{Q}_7^*)} + \frac{\sigma_7^7 \mathrm{Q}_7^{*7}}{(1 + \mathrm{Q}_7^*)^7}$$

$$\frac{4\, b_7 \sigma_7 (\delta_7 - \mathrm{b})\mathrm{Q}^*}{\delta_7 (1 + \mathrm{Q}_7^*)^7} - \frac{4\,(\delta_7 - b_7)^7 (1 - \omega_7)\mathrm{Q}_7^*}{\delta_7 \left(1 + \omega_7 \mathrm{Q}_7^*\right)(1 + \mathrm{Q}_7^*)}\mathrm{P} * (0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_7\, ae_7^{mu}\ [\mu_0]_7 - \frac{a_7}{a_7+7}$$

Hence if $[\mathrm{a}_7^7 > 4(a_7)_7\mathrm{e}$ and

In the following are given some more modular structures that are applied with the model and the concomitant results derived thereof:

Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments

(1) We have an intuitive understanding of how the world works. Our intuition works well in our everyday lives, but it doesn't quite match (e&eb) reality.

(2) One of the basic assumptions authors make about objects is that they are (=) objectively real.

(3) For example, when you leave your apartment in the morning, you imagine your house is right where you left it as you go about your day. Barring some calamity, you expect to find your apartment just as you left it when (e) you return in the evening.

(4) This idea is known as Macrorealism, in that a macroscopic object such as your house is still there even when (e) you're not looking at it

(5) In scientific terms we would say the state of an object exists independent of (e) an observer. While this seems obviously true, we can test the idea experimentally.

(6) To show this, let's consider an experiment. Suppose we have a box that contains a coin. At any time we can open the box and (eb) see whether it's heads or tails.

(7) The catch is (=) that the box has to follow a certain rule. Whatever you observe initially, (eb) you must observe the same thing ten seconds later.

(8) If you open the box and see heads, then (eb) ten seconds later when you open the box you must also see heads.

$$[a_7] = \frac{1K_7 b_7}{\delta_7 - b_7} + \frac{\sigma_7 \, \mathrm{Q}_7^*}{1 + \mathrm{Q}_7^*} > 1\vartriangle$$

This equation is important in its consequence in respect of following cases for which the model is applied:

**Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments**

(7) For the experiment we'll open the box, then open it five seconds later, then finally open it again at ten seconds. We know the first and last observations have to agree, but what about the observation in the middle? It depends upon what the box does with (e&eb) the coin.

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| (8) One possibility is that the box does nothing, so if you see tails the first time (eb), you will see tails both other times. |
| (9) Another possibility is that each time you open the box (eb) it flips the coin when you close it. |
| (10) So the first time you see heads, (eb) the second time tails, and (eb) the third time heads again. |
| (11) But they could also make a tricky box that flips (e&eb) the coin at seemingly random times, but does so in such a way that every ten seconds it repeats (eb) the pattern. In that case, authors would always see the same thing the first and last time, but the middle observation would seem (eb) random. |
| (12)If Macrorealism is true, then (eb) middle observations being always the same as (=) the first and last observation, always opposite or always random are (=) the only possibilities for the experiment |

**Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry Download as PDF (Size: 62KB) HTML PP. 50-53 DOI: 10.4236/jqis.2011.12007 5,508 Downloads 14,712 Views Citations Author(s) Leave a comment Mohamed S. El Naschie**

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| (7) Building upon the pioneering work of J. Bell [1] and an incredible result due to L. Hardy [2] it was shown (eb) that the probability of quantum entanglement of two particles is (=) a maximum of 9.0169945 percent [2]. |
| (8) This happens to be exactly the golden mean to (e) the power of five (?5) [3-7]. |
| (9) Although it has gone largely unnoticed for a long time, this result was essentially established independently in a much wider context by (e) the present author almost two decades ago [3-6]. |
| (10) The present work gives two fundamentally different derivations of (e) Hardy's beautiful result leading to (eb) precisely the same general conclusion, namely that by virtue of (e) the zero measure of the underlying |
| (11)Cantorian-fractal spacetime geometry the notion of (e) spatial separability in quantum physics is (=) devoid of any meaning [7]. The first derivation is purely logical and uses (e) a probability theory which combines the discrete with (e&eb) the continuum. |
| (12)The second derivation is purely geometrical and topological using (e) the fundamental equations of a theory developed by the author and his collaborators frequently referred to as E-infinity or Cantorian spacetime theory [3-7]. KEYWORDS Hardy's Quantum Entanglement, Golden mean, Cantor sets, Fractal Spacetime, E-Infinity Theory, Quantum Mechanics, J. S. Bell Cite this paper M. El Naschie, "Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry," Journal of Quantum Information Science, Vol. 1 No. 2, 2011, pp. 50-53. Doi: 10.4236/jqis.2011.12007. |

**Chaos, Solitons & Fractals Volume 28, Issue 4, May 2006, Pages 857–878 Sir Hermann Bondi 1919–2005 Cantorian spacetime and Hilbert space: Part I—Foundations G. Iovane http://dx.doi.org/10.1016/j.chaos.2005.08.074**

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| (11) Authors are going to show the link between the ε(∞) Cantorian space and( e&eb) the Hilbert spaces H(∞). |
| (12) In particular, El Naschie's ε (∞) is (=) a physical spacetime, i.e. an infinite dimensional fractal space, where (e) time is spacialized and the transfinite nature manifests (eb) itself. |
| (13) El Naschie's Cantorian spacetime is an arena where (e) the physics laws appear at each scale in a self-similar way linked to (e&eb) the resolution of the act of observation. |
| (14) By contrast the Hilbert space H (∞) is (=) a mathematical support, which describes (eb) the interaction between the observer and (e&eb) the dynamical system under (e&eb) measurement. |

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| (15) The present formulation, which is based on (e) the non-classical Cantorian geometry and topology of spacetime, automatically solves (e) the paradoxical outcome of (e) the two-slit experiment and the so-called particle–wave duality. |
| (16) In particular, measurement (i.e. the observation) is equivalent to (=) a projection of ε(∞) in the Hilbert space built on 3 + 1 Euclidean spacetime. |
| (17) Consequently, the wave–particle duality becomes (=) a mere natural consequence of (e) conducting an experiment in (eb) a spacetime with (e&eb) non-classical topological and (e&eb) geometrical structures, while (e) observing and taking measurements in (eb) a classical smooth 3 + 1 Euclidean spacetime. |
| (18) In other words, the experimental fact that a wave–particle duality exists is (=) an indirect confirmation of the existence of ε(∞) and a property of (e) the quantum–classical interface. |
| (19) Another direct consequence of the fact that real spacetime is) =) infinite dimensional hierarchical ε(∞) is (=) the existence of scaling law R(N), introduced by the author, which generalizes the Compton wavelength. |
| (20) It gives an answer to the problem of (e) segregation of matter at different scales, and shows (eb) the role of fundamental constants such as (=) the speed of light and Plank's constant h in the fundamental lengths scale without (e) invoking the principles of quantum mechanics. Copyright © 2005 Elsevier Ltd. All rights reserved |

**Chaos, Solitons & Fractals Volume 28, Issue 4, May 2006, Pages 857–878 Sir Hermann Bondi 1919–2005 Cantorian spacetime and Hilbert space: Part I—Foundations G. Iovane http://dx.doi.org/10.1016/j.chaos.2005.08.074**

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**Chaos, Solitons & Fractals Volume 14, Issue 9, December 2002, Pages 1341–1351 Fractal strings as an alternative justification for El Naschie's Cantorian spacetime and the fine structure constant Carlos Castro, http://dx.doi.org/10.1016/S0960-0779(02)00110-8**

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| (9) Beginning with the most general fractal strings/sprays construction recently expounded in the book by Lapidus and Frankenhuysen, it is shown how the complexified extension of El Naschie's cantorian-fractal spacetime model belongs to (e&eb) a very special class of families of fractal strings/sprays whose scaling ratios are given by (e) suitable pinary (pinary, p prime) powers of the Golden Mean. |
| (10) They then proceed to show why the logarithmic periodicity laws in (eb) Nature are direct physical consequences of (e) the complex dimensions associated with (e&eb) these fractal strings/sprays. |
| (11) Authors proceed with a discussion on quasi-crystals with (e&eb) p-adic internal symmetries, von Neumann's Continuous Geometry, the role of wild topology in (e&eb) fractal strings/sprays, (e&eb) the Banach-Tarski paradox, and tesselations of (e) the hyperbolic plane, quark confinement and (e&eb) the Mersenne-prime hierarchy of (e) bit-string physics in determining the fundamental physical constants in Nature. Fax: +1-914-965-3397 Copyright © 2002 Published by Elsevier Ltd. About ScienceDirectRemote access Shopping cart Contact and support Terms and conditions Privacy policy Cookies are used by this site. For more information, visit the cookies page. Copyright © 2016 Elsevier B.V. or its licensors or contributors. Science Direct ® is a registered trademark of Elsevier B.V. |

**Chaos, Solitons & Fractals Volume 21, Issue 2, July 2004, Pages 261–271 On Cantorian spacetime over number systems with division by zero Jakub Czajko, http://dx.doi.org/10.1016/j.chaos.2003.12.046**

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| (12) Division by zero is proposed in order to open (eb) new horizons for (e) abstract mathematics and physics. |
| (13) Although the formerly forbidden operation may be of (e) little practical value, the feasibility of unrestricted operations is needed for (e) meaningful deployment of linear vector spaces, whose duals are (=) inverse transposes of their primary vector spaces. |
| (14) Both quantum mechanics and (e&eb) pan-geometry of multispatial hyperspace operate on (e&eb) vectors that belong to (e&eb) mutually dual linear vector spaces. |
| (15) The division by zero also may remove (e) many years old paradox related to (e&eb) multiplication of zeros and makes it possible to introduce (eb) the concept of bigroup as (=) an additive and multiplicative group at the same time. |
| (16) This operational extension to (e&eb) algebraic groups makes thus (eb) an infinite-dimensional, fractal Cantorian spacetime (=) the preferred, invertible, operationally unrestricted, abstract mathematical infrastructure for physics. Tel.: +1-925-681-2058 Copyright © 2003 Elsevier Science Ltd. All rights reserved. |

**Chaos, Solitons & Fractals Volume 21, Issue 2, July 2004, Pages 261–271 On Cantorian spacetime over number systems with division by zero Jakub Czajko, http://dx.doi.org/10.1016/j.chaos.2003.12.046**

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The nontrivial equilibrium is stable.

Some of the examples in which stability analysis is of seminal importance are: These are the system to which the

model becomes applicable

**Observation of a Dissipation-Induced Classical to Quantum Transition J. Raftery, D. Sadri, S. Schmidt, H. E. Türeci, and A. A. Houck Phys. Rev. X 4, 031043 – Published 8 September 2014**

(195)Here, authors report the experimental observation of (e) a dynamical quantum phase transition in (eb) a strongly interacting open photonic system.

(196) The system studied, comprising a Jaynes-Cummings dimer realized on (eb) a superconducting circuit platform, exhibits (eb) a dissipation-driven localization transition.

(197)Signatures of the transition in the homodyne signal and photon number reveal (eb) this transition to be from (e) a regime of classical oscillations into (e&eb) a macroscopically self-trapped state manifesting revivals, (=) a fundamentally quantum phenomenon.

(198) This experiment also demonstrates (eb) a small-scale realization of a new class of quantum simulator, whose well-controlled (e&eb) coherent and dissipative dynamics is suited to (e) the study of quantum many-body phenomena out of (e) equilibrium. Received 10 December 2013 DOI:http://dx.doi.org/10.1103/PhysRevX.4.031043 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**Viewpoint: Fermionic Vortices Find their Dual John McGreevy, Department of Physics, University of California at San Diego, La Jolla, CA 92093-0319, USA June 27, 2016• Physics 9, 72Theoretical work reveals a surprising relationship between the physics of fermionic vortices and quantum electrodynamics. Figure caption expand figure APS/Alan Stonebraker Figure 1: The action of discrete symmetries on the fermionic particle (top) and its dual fermionic vortex (bottom).**

(199)Dirac fermions in 2+1 dimensions (two spatial dimensions plus time) have been found in (eb) graphene and on (e&eb) the surfaces of topological insulators (TIs).

(200) A trio of papers now shows that these particles can also be found in (eb) a dual description of quantum electrodynamics (QED) [1–3].

(201)This deceptively simple statement joins (e&eb) together several disparate lines of progress in theoretical physics: interacting topological insulators, (e&eb) anomalies in quantum field theory, (e&eb) spin liquids, and (e&eb) compressible Quantum Hall States.

(202)Humans are not very good at understanding interacting many-body systems.

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| (203)We understand the behavior of free particles and can sometimes use (e) perturbation theory to (e) incorporate weak interactions between (e&eb) them. |
| (204)But what happens when we couple (e&eb) the particles strongly? |
| (205)To reveal the interesting collective phenomena that can result, we've had to rely largely on (e) experiments. |
| (206)Examples include (e) superconductivity (e&eb) superfluidity (e&eb) the integer and (e&eb) fractional quantum Hall effects (e&eb) the fractionalization of charge and spin, and (e&eb) the wetness of water. |
| (207)The fractional quantum Hall effect (FQHE) is a physical phenomenon in which the Hall conductance of 2D electrons shows (eb) precisely quantised plateaus at (eb) fractional values of (e) {\displaystyle e^{2}/h} e^{2}/h. |
| (208) It is a property of (e) a collective state in which electrons bind (e&eb) magnetic flux lines to make (eb) new quasiparticles, and excitations have (e) a fractional elementary charge and possibly also fractional statistics. The 1998 Nobel Prize in Physics was awarded to Robert Laughlin, Horst Störmer, and Daniel Tsui for the discovery and explanation of the fractional Hall effect.[1] |
| (209) However, Laughlin's explanation was (=) a phenomenological guess and only applies to (e&eb) fillings {\displaystyle \nu =1/m} \nu =1/m where (e) {\displaystyle m} m is an odd integer. The microscopic origin of the FQHE remains unknown and is a major research topic in condensed matter physics. (Wikipedia) |
| (210)The quantum Hall effect (or integer quantum Hall effect) is a quantum-mechanical version of (e) the Hall effect, observed in (eb) two-dimensional electron systems subjected to (e&eb) low temperatures and strong magnetic fields, in which the Hall conductance σ undergoes (e&eb) certain quantum Hall transitions to take on (e&eb) the quantized values {\displaystyle \sigma ={\frac {I_{\text{channel}}}{V_{\text{Hall}}}}=\nu \;{\frac {e^{2}}{h}},} \sigma = \frac{I_\text{channel}}{V_\text{Hall}} = \nu \; \frac{e^2}{h}, where {\displaystyle I_{\text{channel}}} I_\text{channel} is the channel current, {\displaystyle V_{\text{Hall}}} V_\text{Hall} is the Hall voltage, e is the elementary charge and h is Planck's constant. The prefactor ν is known as the "filling factor", and can take on either integer (ν = 1, 2, 3, ...) or fractional (ν = 1/3, 2/5, 3/7, 2/3, 3/5, 1/5, 2/9, 3/13, 5/2, 12/5, ...) values. |
| (211) The quantum Hall Effect is referred to as (=) the integer or fractional quantum Hall effect depending on whether ν is an integer or fraction, respectively. |
| (212)The striking feature of the integer quantum Hall Effect is (=) the persistence of the quantization (i.e. the Hall plateau) as (e) the electron density is varied. |
| (213) Since the latter remains constant when (e) the Fermi level is in a clean spectral gap, it must be that the Fermi level is immersed in (eb) essential spectrum (i.e. of finite density of states), though Anderson localized. |
| (214) The quantization and stability of the Hall conductance in such conditions was achieved in 1994 by (e) Bellissard, van Elst and Schulz-Baldes using the methods of (e) Non-Commutative Geometry.[1] |
| (215)The fractional quantum Hall effect is more complicated, as its existence relies fundamentally on (e) electron–electron interactions. |
| (216)Although the microscopic origins of the fractional quantum Hall effect are unknown, there are several phenomenological approaches that provide (eb) accurate approximations. |
| (217) For example the effect can be thought of as (=) an integer quantum Hall effect, not of electrons but of (e) charge-flux composites known as (=) composite fermions. |

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| (218) In 1988, it was proposed that there was quantum Hall effect without Landau levels. This quantum Hall Effect is referred to as (=) the quantum anomalous Hall (QAH) effect. |
| (219)There is also a new concept of the quantum spin Hall effect which is an analogue of the quantum Hall effect, where (e) spin currents flow instead of (e&eb) charge currents.[2] |
| (220)In two dimensions, when classical electrons are subjected to (e&eb) a magnetic field they follow (eb) circular cyclotron orbits. |
| (221)When the system is treated quantum mechanically (eb) these orbits are quantized. |
| (222)The energy levels of these quantized orbitals take on (e) discrete values |
| (223):{\displaystyle \textstyle E_{n}=\hbar \omega _{c}(n+1/2),} \textstyle E_{n}=\hbar \omega _{c}(n+1/2), where ωc = eB/m is the cyclotron frequency |
| (224) These orbitals are known as Landau levels, and at weak magnetic fields, their existence gives rise to (eb) many interesting "quantum oscillations" such as the Shubnikov–de Haas oscillations and the de Haas–van Alphen effect (which is often used to map the Fermi surface of metals). |
| (225) For strong magnetic fields, each Landau level is (=) highly degenerate (i.e. there are many single particle states which have the same energy En). |
| (226)Specifically, for a sample of area A, in magnetic field B, the degeneracy of each Landau level is (=) {\displaystyle \textstyle N=g_{s} BA/\phi _ {0},} \textstyle N=g_{s} BA/\phi _ {0}, where gs represents (eb) a factor of 2 for spin degeneracy, and ϕ0 = 2·10−15 Wb is the magnetic flux quantum. |
| (227) For sufficiently strong B-fields, each Landau level may have (e) so many states that all of the free electrons in the system sit in only a few Landau levels; it is in this regime where (e) one observes the quantum Hall Effect. |
| (228)The integers that appear in the Hall Effect are (=) examples of **topological quantum numbers**. |
| (229) They are known in mathematics as the first **Chern numbers** and are closely related to (e&eb) **Berry's phase.** |
| (230)A striking model of much interest in this context is the Azbel-Harper-Hofstadter model whose quantum phase diagram is (=) the Hofstadter butterfly shown in the figure. (deleted due to spatial constraints) |
| (231)The vertical axis is (=) the strength of the magnetic field and the horizontal axis is the chemical potential, which fixes the electron density. The colors represent the integer Hall conductance. Warm colors represent positive integers and cold colors negative integers. The phase diagram is fractal and has structure on all scales. In the figure there is an obvious self-similarity. |
| (232)Concerning physical mechanisms, impurities and/or particular states (e.g., edge currents) are important for (e) both the 'integer' and 'fractional' effects. |
| (233) In addition, Coulomb interaction is also essential in (eb) the fractional quantum Hall effect. |
| (234)The observed strong similarity between integer and (=) fractional quantum Hall effects is explained by (e) the tendency of electrons to form bound states with (e&eb) an even number of magnetic flux quanta, called composite fermions. |
| (235)Duality is the circumstance in which cranking up interactions results in (eb) another system we understand, usually one that can be described in terms of (e&eb) another set of weakly coupled particles. |
| (236)When such a dual representation is found, (eb) it gives a solution to the strong coupling problem. |
| (237) A further reason that theoretical physicists get excited about dualities is that they undermine (e) the very notion of an “elementary particle.” |
| (238)In terms of the original variables, the dual, weakly coupled particles are (=) generally solitons—large, heavy, classical objects made up of (e) many quanta. |

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| (239)The remarkable thing is that solitons can become (=) "elementary," in the sense that the model that describes the system can be reformulated with the solitons as (=) the basic constituents. |
| (240)An example of a duality, which we will use below, arises in (eb) Maxwell's theory of electromagnetism, where (e) the strength of the interactions is determined by (e) the electric charge e. |
| (241)The duality exchanges electric and (e&eb) magnetic point charges, and the interaction strength in the dual description is 1⁄e, which is small when (e) e is large. |
| (242)This idea has been generalized to (e) many other field theories, thanks, in particular, to tools of (e) supersymmetry [4]. |
| (243)Another classic example of duality involves (e&eb) vortices. |
| (244) A vortex is a localized object in the plane around which a phase field winds (e&eb). (A phase field takes values on the circle.) |
| (245)Dualities relating particles and vortices have (e) a rich history in particle physics and condensed-matter physics. |
| (246) The simplest version relates the particles in the normal state of a superfluid to (e&eb) the vortices in a dual superconductor [5]. |
| (247)This charge-vortex duality has been (=) a workhorse in the study of two-dimensional condensed-matter systems, such as (=) fractional Quantum Hall States, superconducting thin films, and beyond-Landau critical phenomena [6]. |
| (248)Vortices can be (=) fermions (as they frequently are, for instance, in type-II superconductors). |
| (249) A dual description of a system with (e&eb) fermionic vortices thus requires (e) a new ingredient, since (e) it must include (e) such a fermionic particle. |
| (250)The new work [1–3] builds on (e&eb) earlier theoretical work [7] to make (eb) an appealing proposal for (e) the required duality: the fermionic vortex of QED in (eb) 2+1 dimensions is (=) a free Dirac fermion. The argument for this proposal builds on our understanding of TIs. |
| (251)TIs are distinguished only by (e) their edge states. |
| (252) The most famous example is a (3+1)-dimensional time-reversal invariant TI, which can host boundary states whose band structure is defined by (e) a single Dirac cone. |
| (253)A single time-reversal-symmetric Dirac cone is not possible without (e) the bulk TI, so its appearance in experiments provides (eb) a signature of the bulk phase. |
| (254) But the Dirac cone is not (e) the only possible signature. |
| (255) Deformations of the boundary conditions can dramatically change (e&eb) the spectrum of the edge states, yet preserve some essential "TI-ness." This essential property is (=) called an anomaly in high-energy theory. |
| (256)Max Metlitski of the University of California, Santa Barbara, and Ashvin Vishwanath of the University of California, Berkeley [1], and independently, Chong Wang and Todadri Senthil at the Massachusetts Institute of Technology, Cambridge [2], use this ambiguity in the edge theory and the electric-magnetic duality of an auxiliary theory of bulk electromagnetism coupled to (e&eb) the TI to argue for the duality. |
| (257)When I said that (3+1)-dimensional QED had a duality that interchanged the electric and magnetic charges, one could have (e) complained that the electron is (=) a fermion, while (e) a magnetic monopole may not be. |
| (258)However, the authors observe that in a TI, a charge-two monopole is (=) also a fermion. |

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| (259) By analyzing the spectrum of electric and magnetic charges, they conclude that the duality exchanges the electron and (e&eb) the charge-two monopole and, moreover, exchanges a TI with (e&eb) time-reversal symmetry with (e&eb) a "chiral TI," which instead has (e) particle-hole symmetry |
| (260). But this is just what is required of (e) a symmetry that exchanges a fermionic particle with (e&eb) a fermionic vortex: time-reversal symmetry maps the electron to itself but maps the vortex to (e&eb) another vortex whose phase winds (e&eb) the other way (see Fig. 1). |
| (261)The exchange of particles and holes, in turn, preserves (eb) the vortex. |
| (262)The bulk duality implies (eb) a duality between edge theories. |
| (263)The freedom to deform the edge of the TI implies (eb) a weak form of equivalence: although the edge theories are (=) topologically the same (in terms of symmetry assignments and (e&eb) anomalies), they need not have (e) identical low-energy physics. |
| (264)One side of the duality is (=) a free, massless Dirac particle, but the long-distance physics of (2+1)-dimensional QED is (=) a long-standing problem, which is important in the study of spin liquids. |
| (265) Many gauge theories in two dimensions develop (eb) a gap in their energy spectrum because of (e) the proliferation of monopoles. |
| (266) In the realization proposed by Metlitski and Vishwanath [1] and by Wang and Senthil [2], this cannot happen because (e) the monopoles carry electric charge, and indeed the stronger notion of duality as (=) an equivalence of low-energy physics may hold. |
| (267)Evidence for the stronger relation comes from (e) David Mross and colleagues at the California Institute of Technology, Pasadena [3], who have derived a lattice version of the duality using (e) a clever but explicit change of variables. |
| (268)Their construction is particularly interesting because (e) it uses (e) a lattice of one-dimensional models which themselves can only occur at the edge of a TI. |
| (269)Authors explicitly derive the duality between a free electronic Dirac cone and (e&eb) quantum electrodynamics in (eb) (2+1) dimensions (QED3) with (e&eb) N=1 fermion flavors. |
| (270)The duality proceeds via (e&eb) an exact, nonlocal mapping from electrons to (e&eb) dual fermions with (e&eb) long-range interactions encoded by (e) an emergent gauge field. |
| (271)This mapping allows (eb) us to construct (eb) parent Hamiltonians for exotic topological-insulator surface phases, derive (eb) the particle-hole-symmetric field theory of (e) a half-filled Landau level, and nontrivially constrain (e) QED3 scaling dimensions. |
| (272)They similarly establish duality between bosonic topological insulator surfaces and (e&eb) N=2 QED3. |

**Black hole thermodynamics with generalized uncertainty principle Li Xiang, X. Q. Wen (Submitted on 6 Jan 2009 (v1), last revised 18 Mar 2009 (this version, v2))**

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| (273) In the standard viewpoint, the temperature of a stationary black hole is proportional to (=) its surface gravity, $T_H=\hbar\kappa/2\pi$. |
| (274)This is a semiclassical result and the quantum gravity effects are not (e) taken into consideration. |
| (275)This Letter explores a unified expression for (e) the black hole temperature in the sense of (e) a generalized uncertainty principle (GUP). |
| (276) Discussion involves a heuristic analysis of (e) a particle which is absorbed by (e) the black hole. |
| (277)Besides a class of static and spherically symmetric black holes, an axially symmetric Kerr-Newman black hole is considered. Different from the existing literature, authors suggest (eb) that the black hole's irreducible mass represent (eb) the characteristic size in (eb) the absorption process. |

(278)The information capacity of a remnant is also discussed by (e) Bousso's D-bound in de Sitter spacetime. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP 0910:046,2009 DOI: 10.1088/1126-6708/2009/10/046 Cite as: arXiv:0901.0603 [gr-qc] (or arXiv:0901.0603v2 [gr-qc] for this version)

IF we define

$$[a_7]_7 = [f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}]\,P^*_7$$

$$[a_7]_7 = \left[f^{\,7}_{7\,*}\right] P^*_7 - \frac{h^7_*}{h_*} P^*(0)$$

$$e = [g^7_*]_7\, Q^*_7$$

$$\eta_7(t) = \left[h_*\right]_7 \beta_7 \left(\theta_7(t)\right) - \mu_7(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\,[a_7]M_7 - [a_7]_7 W_7 - [n(t)]M_7$$

$$w_7(t) = eM_7$$

Write

$$A = \begin{pmatrix} -[a_7]_7 - [a_7]_7 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)= \begin{pmatrix} [\eta(t)]_7 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_7 \\ W_7 \end{pmatrix}$$

The solution of (32)is

$$X_7(t) = e^{A,t}[X]_7(0) + \int_0 e^{At}$$

## PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION

## ATTRITION MODELS

## MODULE EIGHT

Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal.

## 16. MODEL FORMULATION

**Model is applied to systems given in the following:**

**Researchers blur the line between classical and quantum physics by connecting chaos and entanglement July 12, 2016 by Sonia Fernandez**

(1036) Experimental link between quantum entanglement (left) and (e&eb) classical chaos (right) found using (e) a small quantum computer. Credit: University of California - Santa Barbara Using a small

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| quantum system consisting of three superconducting qubits, researchers at UC Santa Barbara and Google have uncovered a link between aspects of classical and (e&eb) quantum physics thought to be unrelated: classical chaos and (e&eb) quantum entanglement. |
| (1037) Their findings suggest (eb) that it would be possible to use (e) controllable quantum systems to investigate certain fundamental aspects of nature. |
| (1038) "It's kind of surprising because (e) chaos is this totally classical concept—there's (=) no idea of chaos in a quantum system," Charles Neill, a researcher in the UCSB Department of Physics and lead author of a paper that appears in Nature Physics. |
| (1039) "Similarly, there's no concept of entanglement within (eb) classical systems. And yet it turns out that chaos and (e&eb) entanglement are really very strongly and clearly related." |
| (1040) Initiated in the 15th century, classical physics generally examines and describes (eb) systems larger than atoms and molecules. |
| (1041) It consists of hundreds of years' worth of study including Newton's laws of motion, electrodynamics, relativity, thermodynamics as well as chaos theory—the field that studies the behavior of (e) highly sensitive and unpredictable systems. |
| (1042) One classic example of chaos theory is the weather, in which a relatively small change in one part of (e) the system is enough to foil predictions—and vacation plans—anywhere on the globe. |
| (1043) At smaller size and length scales in nature, however, such as those involving (e&eb) atoms and photons and their behaviors, classical physics falls short. |
| (1044) In the early 20th century quantum physics emerged, with its seemingly counterintuitive and sometimes controversial science, including (e) the notions of superposition (the theory that a particle can be located in (eb) several places at once) and entanglement (particles that are deeply linked behave as (=) such despite physical distance from one another). |
| (1045) And so began the continuing search for connections between the two fields. All systems are fundamentally quantum systems, according Neill, but the means of describing in a quantum sense the chaotic behavior of, (e) say, air molecules in an evacuated room, remains limited. |
| (1046) Imagine taking a balloon full of air molecules, somehow tagging them so you could see them and then releasing them (eb) into a room with no air molecules, noted co-author and UCSB/Google researcher Pedram Roushan. |
| (1047) One possible outcome is that the air molecules remain clumped together (e&eb) in a little cloud following the same trajectory around the room |
| (1048) And yet, he continued, as we can probably intuit, the molecules will more likely take off in (eb) a variety of velocities and directions, bouncing off walls and interacting with (e&eb) each other, resting after the room is (=) sufficiently saturated with them. |
| (1049) "The underlying physics is (=) chaos, essentially," he said. |
| (1050) The molecules coming to rest—at least on (eb) the macroscopic level—is the result of (e) thermalization, or of reaching equilibrium after they have achieved (eb) uniform saturation within the system. |
| (1051) But in the infinitesimal world of quantum physics, there is (=) still little to describe that behavior. |
| (1052) The mathematics of quantum mechanics, Roushan said, do not allow for (e) the chaos described by Newtonian laws of motion. |
| (1053) To investigate, the researchers devised an experiment using (e) three quantum bits, the basic computational units of (e) the quantum computer. |

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| (1054) Unlike classical computer bits, which utilize (e) a binary system of two possible states (e.g., zero/one), a qubit can also use (e) a superposition of both states (zero and one) as a single state. |
| (1055) Additionally, multiple qubits can entangle, or link so closely that their measurements will (=) automatically correlate. |
| (1056) By manipulating these qubits with electronic pulses, Neill caused them to interact, rotate and evolve (e&eb, e&eb, eb) in the quantum analog of (e) a highly sensitive classical system. |
| (1057) The result is a map of entanglement entropy of (e) a qubit that, over time, comes to strongly resemble that of (e) classical dynamics—the regions of entanglement in (eb) the quantum map resemble the regions of (e) chaos on the classical map. |
| (1058) The islands of low entanglement in (eb) the quantum map are located in the places of (e) low chaos on the classical map. |
| (1059) "There's a very clear connection between entanglement and (e&eb) chaos in these two pictures," said Neill. |
| (1060) "And, it turns out that thermalization is (=) the thing that connects chaos and (e&eb) entanglement. |
| (1061) It turns out that they are actually the driving forces behind (eb) thermalization. |
| (1062) "What we realize is that in almost any quantum system, including on quantum computers, if you just let it evolve and you start to study what happens as a function of time, (eb) it's going to thermalize," added Neill, referring to the quantum-level equilibration. "And this really ties together the intuition between classical thermalization and (e&eb) chaos and how it occurs (eb) in quantum systems that entangle." |
| (1063) The study's findings have fundamental implications for (e) quantum computing. |
| (1064) At the level of three qubits, the computation is relatively simple, said Roushan, but as researchers push to build increasingly sophisticated and powerful quantum computers that incorporate (e) more qubits to study highly complex problems that are beyond the ability of classical computing—such as those in the realms of machine learning, artificial intelligence, fluid dynamics or chemistry—a quantum processor optimized for such calculations will be a very powerful tool. |
| (1065) "It means we can study things that are completely impossible to study right now, once we get to bigger systems," said Neill.More information: Nature Physics, DOI: 10.1038/nphys3830 , http://www.nature.com/nphys/journal/vaop/ncurrent/full/nphys3830.html Journal reference: Nature Physics search and more info website Provided by: University of California - Santa Barbara |

**Title: coherent control of macroscopic quantum states in a single-Cooper-pair box Authors: Nakamura, Y.; Pashkin, Yu. A.; Tsai, J. S. Publication: Nature, Volume 398, Issue 6730, pp. 786-788 (1999). (Nature Homepage) Publication Date: 04/1999 Origin: NATURE DOI: 10.1038/19718 Bibliographic Code: 1999Natur.398...786N**

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| (1066) A nanometre-scale superconducting electrode connected to (e&eb) a reservoir via (e&eb) a Josephson junction constitutes (e) an artificial two-level electronic system: a single-Cooper-pair box. |
| (1067) The two levels consist of (e) charge states (differing by 2e, where (e) e is the electronic charge) that are coupled by (e) tunnelling of Cooper pairs through (e&eb) the junction. |
| (1068) Although the two-level system is macroscopic, containing (e) a large number of electrons, the two charge states can be (=) coherently superposed. |

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| (1069) | The Cooper-pair box has therefore been suggested as (=) a candidate for a quantum bit or `qubit'- the basic component of a quantum computer. |
| (1070) | Here authors report the observation of quantum oscillations in (eb) a single-Cooper-pair box. |
| (1071) | By applying a short voltage pulse via (e&eb) a gate electrode, we can control (e&eb) the coherent quantum state evolution: the pulse modifies (e&eb) the energies of (e) the two charge states non-adiabatically, bringing them into (e&eb) resonance. |
| (1072) | The resulting state-a superposition of (e) the two charge states-is detected by (e) a tunnelling current through (e&eb) a probe junction. |
| (1073) | Results demonstrate electrical coherent control of (e) a qubit in a solid-state electronic device. |

**Superconducting quantum computing From Wikipedia, the free encyclopedia**

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| (1074) | Superconducting quantum computing is a promising implementation of (e) quantum information technology that involves (e&eb) nanofabricated superconducting electrodes coupled through (e&eb) Josephson junctions. |
| (1075) | As in a superconducting electrode, the phase and (e&eb) the charge are conjugate variables. |
| (1076) | There exist three families of superconducting qubits, depending on (=) whether the charge, (e&eb) the phase, or neither of (e) the two is good quantum number. |
| (1077) | These are respectively termed charge qubits (e&eb) flux qubits, and (e&eb) hybrid qubits. |

**Theory**

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| (1078) | Unlike many other physical implementations of a qubit which involve (e&eb) exclusively two-level systems (such as nuclear spin and (e&eb) photon polarization), the integrated quantum circuit involved in (e&eb) a superconducting qubit is (=) a multi-level system, of which only the first two levels are used as (=) the computational basis. |
| (1079) | A basic requirement for such an implementation is that the energy levels are (=) not uniformly spaced, so that (e) photons of a particular frequency which cause (eb) transition between the 0 and 1 levels do not(e) cause transitions from the first level to (e&eb) the higher levels as well. |
| (1080) | For electronic signals to be carried from one part of the circuit to (e&eb) another without energy loss and (e&eb) hence decoherence, the system also needs to be (=) non-dissipative, i.e., the metallic parts involved should have (e) zero resistance. |
| (1081) | These circuits currently need to be (=) operated at very low temperatures so that, 1) superconductivity is (=) realized, and 2) thermal fluctuations do not (e) cause transitions between (e&eb) energy levels. |

**Quantum LC Circuit**

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| (1082) | The simplest non-dissipative quantum circuit consists simply of (e) an inductor L {\displaystyle L} L and (e&eb) capacitor C {\displaystyle C} C, with (e&eb) the metallic wires connecting them being superconducting. |
| (1083) | Here the flux Φ {\displaystyle \Phi } \Phi through (e&eb) the inductor and (e&eb) the charge Q {\displaystyle Q} Q in the capacitor are (=) canonically conjugate variables which obey (eb) the commutation relation [ Φ , Q ] = i ℏ {\displaystyle [\Phi ,Q]=i\hbar } [\Phi ,Q]=i\hbar |

**The Hamiltonian associated with this system is**

(1084) H = Φ 2 2 L + Q 2 2 C {\displaystyle H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}} H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}

(1085) This gives rise to the energy levels E = ℏ ω 0 ( n + 1 2 ) {\displaystyle E=\hbar \omega _{0}(n+{\frac {1}{2}})} E=\hbar \omega _{0}(n+{\frac {1}{2}}) Here ω 0 = 1 L C {\displaystyle \omega _{0}={\frac {1}{\sqrt {LC}}}} \omega _{0}={\frac {1}{{\sqrt {LC}}}}

**Quantum gravity Carlo Rovelli (2008), Scholarpedia, 3(5):7117. doi:10.4249/scholarpedia.7117 revision #135539 [link to/cite this article] Post-publication activity Curator: Carlo Rovelli Dr. Carlo Rovelli, Center of Theoretical Physics, and Marseille, France**

(1086) Quantum Gravity is the name given to any theory that describes gravity in (eb) the regimes where quantum effects cannot be disregarded.

(1087) At present, there is no such a theory which is universally accepted and confirmed by experience. Therefore the term "Quantum Gravity" indicates (eb) more an open problem than a specific theory.

(1088) Several research lines, at different levels of development, offer tentative solutions to the problem. These tentative quantum-gravity theories are variously viewed as competing research directions, or as contributions to the common goal of finding the physically correct theory. The quest for the good quantum theory of gravity bears on a number of fundamental issues and it is sometimes presented as the most important open problem in fundamental physics: the "Holy Grail" of contemporary theoretical physics.

**The problem of quantum gravity**

(1089) The current theory of (e) gravity is general relativity, a theory developed by Albert Einstein at the beginning of the XX century, which has changed in depth our understanding of (e) what are space and time.

(1090) Similarly, quantum mechanics has changed in depth our understanding of matter, energy and causality. These two theories, today very well empirically supported, have opened (eb) a major conceptual revolution in physics.

(1091) However, they appear to be incompatible, at least at a first reading, because (e) each of the two is formulated on the basis of (e) principles that are explicitly contradicted by (e) the other theory.

(1092) Therefore in spite of the immense development of scientific knowledge it has delivered, the physics of the XX century has left us with (e&eb) a major confusion on what are the basic conceptual ingredients for (e) understanding the physical world.

(1093) In other words, the scientific revolution opened by general relativity and quantum mechanics at the beginning of (e) the XX century is not concluded yet, and a new synthesis is required. Quantum gravity, merging general relativity and quantum mechanics should be this new synthesis.

**Quantum space and quantum time**

(1094) Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.

(1095) In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.

(1096) They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.

(1097) When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.

(1098) Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.

**The Planck scale**

(1099) Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".

(1100) Here $\hbar$ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.

(1101) The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.

(1102) Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)

(1103) Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.

## Physical applications

Some phenomena where quantum gravity plays a major role are the following.

**The microscopic structure of spacetime**

(1104) If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.

(1105) Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length

(1106) Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.

(1107) A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.

**Early cosmology**

(1108) According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.

(1109) However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.

(1110) A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.

(1111) Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".

(1112) One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.

**Black holes**

(1113) Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.

(1114) First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.

(1115) Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.

(1116) What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.

(1117) Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.

**Astrophysical effects**

(1118) Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.

(1119) An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.

(1120) The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.

**Main approaches**

(1121) The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.

(1122) Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.

(1123) The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.

**String theory**

(1124) String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.

(1125) Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.

(1126) String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.

(1127) Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.

(1128) However, when summed the entire series appears to be (=) divergent.

(1129) In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.

(1130) Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.

(1131) Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.

(1132) In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.

**Loop quantum gravity and spin foams**

(1133) Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.

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| (1134) The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity. |
| (1135) Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime. |
| (1136) The granular properties of space can be (=) explicitly computed. |
| (1137) In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory. |
| (1138) The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks. |
| (1139) Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure. |
| (1140) The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit. |
| (1141) Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism. |
| (1142) This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism. |

**Noncommutative geometry**

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| (1143) Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime. |
| (1144) Quantum mechanics teaches us that dynamical quantities are (=) noncommutative --meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed. |
| (1145) It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry. |
| (1146) A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces. |
| (1147) Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces. |
| (1148) This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity. |

**Others**

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| (1149) Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays |

(1150) . Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection.

(1151) Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible.

(1152) The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.

Let $p_8(t, a_8)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(g_8)Q_8$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_8(0) = 0, \quad g_8(P_8) > 0, \quad P_8 \; 1 0, \qquad \lim_{P_8 \to \infty} g_8(P_8) = (K_8)_8 \; ?\infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_8 + (\rho_a)_8 + D_8(a_8, P_8, Q_8)\rho_8 + 0, \qquad 0 < a_8 \; ?\infty,$$

$$\rho_8(t,0) = \int_0^\infty n_8(a_8, P_8, Q_8)\rho_8(t_8, a_8)da_8$$

Where $D_8$ is the mortality function and $n_8$ is the fertility function.

Put $_8$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_8}{dt}\right]_8 = x_8\left[\frac{dQ_8}{dt}\right]_8 - x_8\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_8}{dt}\right]_8 = x_3\left[\frac{dQ_3}{dt}\right]_8 - x_4\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_8}{dt}\right]_8 - x_6\left[\frac{dQ_3}{dt}\right]_8$$

Where $x_8, x_8, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_8$ and $g_8(P_8)Q_8$

Now when you add the term $(b_8Q_8)$ with $g_8(P_8)(Q_8)$ we are infect see is that $(b_8Q_8)$ is credited incrementally part of $g_8(P_8)(Q_8)$ or instant nearly, when subtracted from $\frac{dQ_8}{dt}$ , it means that $g_8(P_8)(Q_8)$ is slowly eaten by $\frac{dQ_8}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_8 = \mu_8(a_8) + (f_8)_8(P_8)_8 + (f_8)_8(Q)_8 \qquad (2)$$

$$n_8 = \beta_8(a_8)h_8(Q_8), \qquad (3)$$

Where $\mu_8$the age- dependent debit rate whit is $\mu_8(a_8)$ 0 $\forall a > 0$; $(f_8)_8$ is the density- dependent death rate with $f_8(0) = 0$, $(f_8^8)_8(P_8)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_8)_8(P_8) = (k_8)_8$ ; $f_8$ is the functional response of creditors to prey with $[f_8]_8(0) = 0, (f_8^8)_8(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_8)_8(Q_8) = [k_8]$ ; $\beta$ is the age

dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^8(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_8 + (\rho_a)_8 = -\mu_8(a_8) + (f_8)_8 P_8) + \big(f_8)_8(Q_8)\big)(\rho_8)_8 \quad (4)$$

$$\rho_8(t_8, 0) = h_8(Q_8)\int_0^\infty \beta_8(a_8)\rho_8(t, a_8)da, \quad (5)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(P_8)Q_8 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_8$ and $Q_8$ of (4-6) satisfy

$$0 \ \rho_8(t, a_8) \ \rho_8^*(t, a_8),$$

$$0 \ Q_8(t) \ Q_8^*(t),$$

Where $\rho_8^*(t, a_8)$ and $Q_8^*(t)$, are solutions of

$$(\rho_t^*)_8 + (\rho_a^*)_8 = -(\mu_8(a_8) + (f_8)_8(P_8^*) + (f_8)_8 Q_8^*))\,\rho_8^*, \quad (7)$$

$$\rho_8^*(t, 0) = \int_0^\infty \beta_8(a_8)\,\rho_8^*(t, a_8)da_8 \quad (8)$$

$$\frac{dQ_8^*}{dt} = -b_8 Q_8^* + g_8(P_8^*)Q_8^*, \quad (9)$$

With $P_8^*(t) = \int_0^\infty \rho_8^*\ (t, a_8)da_8$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_8^*(t)}{dt} = P_8^*(t)\big(r_8 - (f_8)(P_8^*) - (f_8)_8(Q^*)\big), \tag{10}$$

$$\frac{dQ_8^*(t)}{dt} = Q_8^*(t)(-b_8 + g_8(P_8^*)), \tag{11}$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_8(a)e^{-\int_0^a \mu_8(\tau)dt}e^{-ra}\ da_8 = 1. \tag{12}$$

If $r < 0$, the trivial solution $P_8^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(t)dt}\ da_8 > 1, \tag{$H1$}$$

Which guarantees $r_8 > 0$.

*Remark.* If $\beta_8(a) = [\beta_0]_8 a_8 e^{-\alpha_8 a_8}, \mu_8(a_8) \equiv \mu_8$ then the solution of (12) is

$$r_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_8]_8(P_8) = 0, [f_8](Q_8) = \lambda_8 Q_8\ h_8(Q_8) = 1, \qquad and\ \ g_8(P_8) = c_8 P_8$$

$$[\rho_8]_8 + [\rho_a]_8 = -(\mu_8 + \lambda_8 Q_8)\rho_8$$

$$\rho_8(t,0) = \int_0^{\infty} [\beta_0]_8 a_8 e^{-\alpha_8 a_8} \rho_8(t,a_8)\, da_8$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + c_8 P_8 Q_8$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_8}{dt} = \omega_8 P_8 - \lambda_8 P_8 Q_8$$

$$\frac{dQ_8}{dt} = b_8 Q_8 - c_8 P_8 Q_8$$

Where $\omega_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$

## 17. THE EQUILIBRIUM $(\rho_0)_8(a_8), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_8(a_8), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_8 + (\rho_a)_8 = -\big(\mu_8(a_8) + (f_8)_8(P_8)\big)\rho,$$

$$\rho_8(t,0) = \int_0^{\infty} \beta_8(a_8)\rho_8(t,a_8)da_8 \qquad (13)$$

The solution of (13) has the form of

$$\rho_8(t,a_8) = Ce_8^{r(t-a_8)} - \int_0^a (\tau)d\tau - \int_0^8 (f_8)_8\big(P_8(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_8(t) = C_8 e^{r_8 t - \int_0^t (f_8)_8\left(P_8(\tau)\right)d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau}\, da.$$

Hence, there exists an equilibrium $P_0 = \int_8^{-8}(r)$ if $k_8 > r$. This is equivalent to

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{(k_8)_8 a_8} da_8 < 1.$$

In fact, if we define

$$R_8(s_8) = \int_0^\infty \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau)d\tau} e^{-s_8 a_8}\, da_8$$

$R(s)$ is a decreasing function. $R_8(r_8) = 1$ And $R_8(k_8)_8 < 1$ imply $r_8 < (k_8)_8$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_8(a_8), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{k_8 a_8} da_8 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_8 \equiv \int_0^\infty (\rho_0)_{88}(a_8)da_8$ *satisfies*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{-(f_8)_8(P_0)a_8} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_8(a_8), 0)$

Let $u_8(t,a_8) = \rho(t,a_8) - \rho_0(a_8)$ Then

$$N_8(t) \equiv \int_0^\infty u_8(t,a_8)da_8 = \int_0^\infty \rho_8(t,a_8)da_8 - \int_0^\infty (\rho_0)_8(a_8)da = P_8(t) - P_0)_8$$

And

$$(u_8)_8 + (u_a)_8 = (\rho_t)_8 + (\rho_a)_8 - (\rho_0^8)_8(a_8)$$

$$= -\mu_8(a_8) + (f_8)_8(P_8) + (f_8)_8(Q_8))\,(u_8 + (\rho_0)_8) + \mu_8(a_8)(\rho_0)_8 + (f_8)_8(P_0)_8(\rho_0)_8$$

$$= -\mu_8(a_8)u_8 - (f_8)_8(P_8) + (f_8)_8(Q_8))\,u_8 - (\rho_0)_8\,(f_8)_8(P_8) - (f_8)_8(P_0)_8$$

$$+ (f_8)_8(Q_8)), \qquad (15)$$

Where $(\rho_0^8)_8(a_8)$ denotes $d(\rho_0)_8(a_8)/da_8$

Expansion $(f_8)_8(P_8)$ and $(f_8)_8(Q_8)$ at $(P_0)_8$ and 0 yield respectively

$$(f_8)_8(P_8) = (f_8)_8((P_0)_8 + N_8)$$

$$= (f_8)_8(P_0)_8 + (f_8^8)_8(P_0)_8N_8 + o_8(N_8), \qquad (16)$$

$$(f_8)_8(Q_8)$$

$$= (f_8^8)_8(0)Q + o_8(Q_8) \qquad (17)$$

Write $f_8(P_0) = f_{80}, f_8'(P_0) = f_{80}'$, and $f_8'(0) = f_{80}'$ and substitute (16) and (17) into (15) to obtain

$$(u_8)_8 + (u_a)_8 \approx -(\mu_8(a_8) + (f_{80})_8)u - (f_{80}^8)_8(\rho_0)_8(N_8) - (f_{80}^8)_8(\rho_0)_8(Q_8)$$

$$u_8(t,0) = \rho_8(t,0)(\rho_0)_8(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a) + (\rho_0)_8(a_8)da_8 - (\rho_0)_8\,(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a_8)da_8 + (h_8(Q_8) - 1)(\rho_0)_8\,(0) \qquad (18)$$

Expanding

$$h_8(Q_8) = 1 + (h^8)_8(0)Q_8 + 0(Q_8) \equiv 1 + h_0^{'}Q_8 + 0Q_8$$

And substituting it into (18), we arrive at

$$u_8(t,0) \approx \int_0^{\infty} \beta_8\,(a_8)u_8(t,a_8)da_8 + (h_{a_8}^{'})(\rho_0)_8 Q_8$$

The linearization is the following :

$$[u_8] + [u_8] = -(\mu_8(a_8) + [t_{80}]_8 u_8 - [\alpha_8]_8(a_8)N_8 - [\alpha_8]_8(a_8)Q_8$$

$$u_8(t,0) = \int_0^{\infty} \beta_8(a_8)u_8(t,a_8)da_8 + [\xi_8]Q_8 \qquad (19)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8$$

where

$$[\alpha_8]_8(a_8) = [t_8^8 a]_8[p,o]_8(a_8)d2(a_8) = [f_{80}^8]_8(\rho\, o\,)_8(a_8)\xi = [\mathrm{h\ o}^8]_8[\rho \mathrm{o}]_8(\mathrm{o}) \text{ and } \mathrm{b}_8 = \mathrm{b}_8 - g_8(P_8)$$

If $g_8(P_0) > b$, clearly, $(\rho_0)_8(a_8), 0)$ is unstable and the predator population is persistent. Assume $g_8(P_0)_8 < b$.

Then, $Q_8(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_8 + u_a = -(\mu_8(a_8) + (f_{80})_8)u_8 - (\alpha_8)_8(a_8)N_8$$

$$= -(\mu_8(a_8) + f_{80})u_8 - \alpha_8(a_8)\int_0^{\infty} u_8(t,a_8)da_8$$

$$u_8(t,0) = \int_0^\infty \beta_8\,(a_8)u_8(t,a_8)da_8$$

Using the method developed by Cushing [2, 3], we suppose

$$u_8(t,a_8) = \tilde{u}_8(a_8)^{e^{\gamma(t-a_8)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_8}\,da_8$. Then,

A solution of the first differential equation for $\tilde{u}(a_8)$ gives,

$$\tilde{u}_8(a_8) = e^{-\int_0^a \mu_8(\tau_8)d\tau - (f_{80})_8 a}\left(\tilde{u}_8(0) - w\int_0^a e^{\int_0^\tau \mu_8(s_8)ds + (f_{80})}\tau_8\,(\alpha_8)_8(\tau_8)e^{\gamma_8\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\right)\tilde{u}_8(0)$$

$$+w\int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\int_0^a (\alpha_8)_8(\tau_8)e^{-\gamma_8\tau_8}e^{(f_{80})_8\tau_8 + \int_0^\tau \mu_8(s_8)ds}d\tau_8\,da_8$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}da\right)(\tilde{u}_8(0))_8$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_8 a_8-(f_{80})_8 a_8-\int_0^a \mu_8(\tau_8)d\tau}\int_0^a e^{(f_{80})_8\tau_8+\int_0^\tau \mu_8(s_8)ds+y\tau}\alpha_8(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_8 a_8-(f_{80})_8 a_8-\int_0^a \mu_8(t_8)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_8\,(a_8)Y_8(\gamma_8,a_8)da_8\right)u_8(0)$$

$$+w\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)\int_0^a(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8=0,$$

(24)

$$\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\tilde{u}_8(0)$$

$$+\left(-1-\int_0^{\infty}Y_8(\gamma_8,a_8)\int_0^a(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8\right)w=0.$$

Equation (20) has a solution of the form of $u_8(t,a_8)=u_8(a_8)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_8(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_8(\gamma_8,a_8)=\int_0^a(\rho_0)_8(\tau_8)Y^{-8}\,(\gamma_8,a_8)d\tau_8$. We obtain the characteristic equation

$$f_{80_8}^{l}\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\left(\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$+\left(1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)da\right)\left(1+(f_{80}^{l})_8\int_0^{\infty}Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_8(a_8) = (\rho_0)_8(0)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - (f_{80})_8 a_8}$$

$$I_8(\gamma_8, a_8) = \int_0^a (\rho_0)_8(\tau_8)e^{\gamma_8\tau_8 + (f_{80})_8\tau + \int_0^\tau \mu_8(s_8)ds_8}d\tau$$

$$= \frac{(\rho_0)_8(0)}{\gamma_8}(e^{\gamma_8 a_8} - 1)$$

Let

$$E_8(a_8) = e^{-(f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau_8}$$

Then $[Y_8]_8[\gamma_8]_8, a_8) = E_8(a_8)e^{-\gamma_8 a_8}$. by a simple calculation,

$$\int_0^\infty \beta_8(a_8)Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma}\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}\ da\right);$$

$$\int_0^\infty Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma_8}\left(\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}\ da_8\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_8(0)}{\gamma_8}[f_{80}^8]\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}da_8\ (\int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8)$$

$$+\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right).$$

$$\left(1+\frac{(\rho_0)_8(0)}{\gamma_8}(f_{80}^{'})_8\left(\int_0^{\infty}E_8(a_8)da_8-\int_0^{\infty}E_8(a_8)e^{-\gamma_8 a_8}da_8\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right)(1+\frac{(f_{80}^{'})_8}{\gamma_8}(\rho_0)_8(0)\int_0^{\infty}E_8(a_8)da_8)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)da_8=1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\neq 0,$$

Since $\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)\int_0^{a}\alpha_8(\tau_8)Y^{-8}(\gamma_8,\tau_8)d\tau_8 da_8>0$. Therefore, the characteristic solution is

$$\gamma_8=-(f_{80}^{'})_8(\rho_0)_8\int_0^{\infty}E_8(a_8)da_8.$$

Since $f_8^{'}(P_0)_8>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_8),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_8),0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_8 + (\mu_8(a_8) + [f_8]_8(P_8^*) + [f_8]_8(Q_8^*))P_8^* = 0, \quad (26)$$

$$p_8^*(0) = h_8(Q_8^*)\int_0^\infty \beta_8(a_8)p_8^*(a_J)p_8^*(a_8)da_8 \quad (27)$$

$$g_8(P_8^*) - b_8 = 0, \quad (28)$$

Where $P_8^* = \int_0^\infty p_8^*(a_8)\, da_8$

Equation (28) yields $P^* = g^{-8}(b)$,

From (26),

$$p_8^*(a) = p_8^*(0)e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}$$

Substituting into (27), we need to solve

$$\left(1 - h_8(Q_8^*)\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}da_8\right)p_8^*(0) = 0$$

For $Q_8^*$

Fix $p_8^*(0)$ and $P_8^*$ Let

$$H_8(z_8) = 1 - h_8(z_8)\int_0^\infty \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}da_{,8}$$

And assume that

$$[k_3]_8 \int_0^{\infty} \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(g_8^{-8}(b_8))a_8 - [k_8]_8 a_8}\, da_8 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty} [H_8(z_8)] > 0.$$

Clearly,

$$\frac{dH_8}{dz_8} = h_8(z)(\,[f_8^{'}]_8(z_8)\,) \int_0^{\infty} a_8 \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - f_8(P_8^*)a_8 - f_8(z_8)a_8} da_8$$

$-h^{'}(z_8) \int_0^{\infty} \beta_8(a_8)\, e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}\ da_8$

By the assumptions on $\left[f_8^8\right]_8 and\ [h^8]^8 dH_8/dz > 0$. Hence, if

$$H_8(0) = 1 - \int_0^{\infty} \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8} da_8 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_8(0) < 0$ follows if $g_8^-(b_8) < [f_8^{-8}]_8(r_8)$, where $r_8$ is the solution of (12).

Indeed, since

$$R_8(s_8) = \int_0^{\infty} \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 e^{-y_8^a}} da_8$$

Is a decreasing function of $s_8$ when $P_8^* = [g^{-8}]_8(b_8) < [f_8^{-8}]_8(r_8)$, that is $[f_8]_8(P_8^*) < r$, $R\big(f_8(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_8(g^{-8}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8\left(g_8^{-8}(b)\right)a_8} > 1. \qquad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_8^*(a_8), Q_8^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_8^*(a_8), Q_8^*)$ is more complicated. We again consider perturbations $m_8(t_8 a_8) = p_8(t_8, a_8) - p_8^*(a_8)$ and $W_8(t) = Q_8(t) - Q^*$.Then

$$M_8(t) \equiv \int_0^\infty m_8(t, a_8)da_8 = P_8(t) - P_8^*,$$

$$[m_8]_8 + [m_a]_8 = [p_8]_8 + [p_a]_8 - \left[p^*\right]_8(a_8) \approx -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8$$

$$-[f_{8*}^8]_8 p_8^*(a_8)M - [f_{8*}^8]_8 p_8^*(\mathrm{a})W_8,$$

$$m_8(t,0)p_8(t,0) - p_8^*(0) \approx h_* \int_0^\infty \beta_8(a_8)m_8(t,a_8)\,da + \frac{[h_*^8]_8}{[h_*]_8}\, p_8^*(0)W_8$$

$$\frac{dW_8}{dt} = \frac{dQ_8}{dt} \approx [g_*^8]_8 Q_8^* M_{,8}$$

Where $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $\left[f_{8*}^{'}\right]_8$ denoting $\left[f^{'}{}_8(P^*)\right]_8\left[f_{8*}^{'}\right]_8$ denoting $[f_8]_8$ $(Q_8^*,$ $h_{*}^{'}$ denoting $h_*^{'}$ denoting $h_8^{'}$ $(Q_8^*)$, and $\left[g_*^{'}\right]_8$

Denoting $\left[\dot{g}\right]_8(P_8^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_8]_8 + [m_a]_8 = -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8 - \left[\dot{f}_{8*}\right]_8[p_8^*](a_8)M_8 - \left[\dot{f_{8*}}\right]_8 p_8^*(a_8)W_8$$

$$m_8(t,0) = [h_*]_8 \int_0^\infty \beta_8(a_8)m_8(t,a_8)da_8 + \frac{[h_*^8]_8}{[h^*]_8}[p^*]_8(0)\, W_{,8} \qquad (30)$$

$$\frac{dW_8}{dt} = [g_*^8]_8[Q_8^*]M_8$$

Integrating $[m_8]_8(t,a_8)$ with respect to a from 0 to ∞, we have

$$\frac{dM_8(t)}{dt} = \int_0^\infty [mt]_8(t,a_8)da_8$$

$$= -\int_0^\infty [m_a]_8\,(t,a_8)da_8 - \int_0^\infty \mu(a_8)m(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + \left[\dot{f}_{8*}\right]_8 \int_0^\infty p_8^*(a_8)da_8\right)M_8 - [f_{8*}^8]_8 \int_0^\infty p_8^*(a_8)da_8 W_8$$

$$= m_8(t,0) - \int_0^\infty \mu_8(a_8)\, m_8(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + \left[\dot{f_{8*}}\right]_8 P_8^*\right)M_8 - \left[\dot{f_{8*}}\right]_8 P_8^* W_{,8}$$

From (30), we arrive at the equations

$$\frac{dW_8(t)}{dt} = \int_0^\infty ([h_*]_8\beta_8(a_8) - \mu_8(a_8))m_8(t,a_8)da_8$$

$$-\left([f_{8*}] + [f_{8*}]_8 + [f_{8*}']_8 P_8^*\right) M_8 + \left(\frac{[h_*']_8}{[h_*]^8}\, P_8^*(0) - f_{8*}' P_8^*\right) W_8,$$

(31)

$$\frac{dW_8(t)}{dt} [g'_*]_8\, Q_8^*\, M_{.8}$$

By the mean value theorem,

$$\int_0^\infty ([h_{"*}]_8 \beta_8(a_8) - \mu_8(a_8)) m_8(t, a_8) da_8 = ([h_*]_8 \beta_8(\theta_8) - \mu_8(\theta_8)) M_{,8}$$

Where $0< \theta_8 = \theta_8(t) < \infty$.

And if a $^8_8$ [a $^8_8$]2<4 [a $_8$]2<4 a $_8$e,

Re $\lambda_8 [\lambda_8]_8 = \text{Re}\lambda_8 [\lambda_8]_8 = -\frac{8}{8} a_8 [\text{a2}\ ]_8 < 0$.

Let ⊿=max$_{8<u<\infty}$ |h $_*$ [h $_*$] $_8$β $_8$($a_8$) − $\mu_8(a_8)$|. Then|$n_8$(t)| ⊿∀t . Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium(ρ $^*$ (a), $\boldsymbol{Q}^*$) is asymptotically stable if either a $^8_8$ <4 a $_8$ e and [a $_8$] > 1⊿,or [a $^8_8$] ⊿[a $_8$]e and (a $_8$ − $\sqrt{(a^8_8)_8 - 4(a_8)_8 e}$ > 1⊿

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+8}\right)$ and β $_8$($a_8$) = ae-$^{ma}$ Then ⊿=[$\mu_0$] + 2 Simple calculation leads to

$$\int_0^{\infty} \beta_8(a_8)\mathrm{e}^{-\int_0^{\mathrm{a}} \mu_8(\tau_8) + d\tau_8 \mathrm{da}} = \frac{1}{(\mathrm{m}+\mu_0+1)^8} + \frac{1}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_8 + [\mu_0]$ 0.5. Assume$[f_8](x_8) = K_8 x_8$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_8(x_8) = [0x_8]/(2 + x_8)$; the functional responses $[[f_8](x) = [0]\, x_8/(1 + x_8)$and $h_8(x_8)$=$(1+\omega_8 x_8)/(1 + x_8)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_8 + [\mu_0] + [\mathrm{K}][P_0])}{[m_8]_8 + (\mu_0)_8 + [\mathrm{K}][P_0])_8^3\ [\mathrm{K}]_8[P_0]_8^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_8$i e $[P_0] < b/(\delta_8 - b_8)$then the equilibrium$(P_0)$ $[P_0]_8(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_8$

$[f_{8*}] = K_8\ \mathrm{P}^*\ [f_{8*}] = \frac{\sigma \mathrm{Q}_8^*}{8+\mathrm{Q}_8^*} \quad [h_*] = \frac{8+\omega_8 \mathrm{Q}_8^*}{8+\mathrm{Q}_8^*}$

$[f_{8*}] = K_8\ [\mathrm{f}_{8*}^8] = \frac{\sigma_8}{(8+\mathrm{Q}_8^*)^8}[\ \mathrm{g}_*^8] = \frac{\sigma_8}{(8+\mathrm{P}_8^*)^8} \quad [\mathrm{h}_*^8] = \frac{\omega_8 - 8}{(8+\mathrm{Q}_8^*)^8}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta$, and $\omega$ need to satisfy

$$\frac{\omega_8}{(m + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 + b_8} + \sigma)^8} + \frac{1\omega_8}{([m_8]_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 + b_8} + \sigma_8)^3} < 1$$

And

$$\frac{1}{(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^8} + \frac{1\omega_8}{(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8} + (m_8 + [\mu_0]_8 + \frac{K_8 b_8}{\delta_8 - b_8} + \sigma_8)\, Q_8^*)^3$

$= m_8 + [\mu_0]_8 + 3 + \frac{K_8 b_8}{\delta_8 - b_8}(1 + Q^*)(1 + \omega_8 Q_8^*$

Now

$$\left[a_{11}\right] - 4\left[a_{88}\right] e = \left(\frac{1K_8 b_8}{\delta_8 - b_8}\right)^8 + \frac{4K_8 b_8 \sigma_8 Q_8^*}{(\delta_8 - b_8)(8 + Q_8^*)} + \frac{\sigma_8^8 Q_8^{*8}}{(1 + Q_8^*)^8}$$

$$\frac{4\, b_8 \sigma_8 (\delta_8 - b) Q^*}{\delta_8 (1 + Q_8^*)^8} - \frac{4\, (\delta_8 - b_8)^8 (1 - \omega_8) Q_8^*}{\delta_8 (1 + \omega_8 Q_8^*)(1 + Q_8^*)} P^*(0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_8\, ae_8^{mu}\ [\mu_0]_8 - \frac{a_8}{a_8 + 8}$$

Hence if $[a_8^8 > 4(a_8)_8 e$ and

$$[a_8] = \frac{1K_8 b_8}{\delta_8 - b_8} + \frac{\sigma_8 Q_8^*}{1 + Q_8^*} > 1\Delta$$

**The nontrivial equilibrium is stable.**

IF we define

$$[a_8]_8 = [f_{8*}]_8 + [f_{8*}]_8 + [f_{8*}] P_8^*$$

$$[a_8]_8 = \left[f_{8*}^{8}\right] P_8^* - \frac{h_*^8}{h_*} P^*(0)$$

$$e = [g_*^8]_8 Q_8^*$$

$$\eta_8(t) = \left[h_*\right]_8 \beta_8 \left(\theta_8(t)\right) - \mu_8(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\left[a_8\right] M_8 - [a_8]_8 W_8 - [n(t)] M_8$$

$$w_8(t) = eM_8$$

Write

$$A = \begin{pmatrix} -[a_8]_8 - [a_8]_8 \\ e \qquad 0 \end{pmatrix}$$

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| $B=(t)=\begin{pmatrix} [\eta(t)]_8 & 0 \\ 0 & 0 \end{pmatrix}$ |
| And |
| $X(t)=\begin{pmatrix} M_8 \\ W_8 \end{pmatrix}$ |
| The solution of (32)is |
| $X_8(t) = e^{A,t}[X]_8(0) + \int_0 e^{At}$ |
| This equation notwithstanding its importance in all the models stand out for its importance for the following models<br><br>to which it is applied: |
| **Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016** |

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| (164)These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field. |
| (165)The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets. |
| (166) Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state. |
| (167) Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands. |
| (168)The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI |
| (169) The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9]. |
| (170)This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive. |

(171)More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

(172)Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.

(173)Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.

(174) The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.

(175)This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.

(176)The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.

(177)This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.

(178) For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.

(179)The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(180)Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.

(181)We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.

(182)When this is true—as is often the case—the system serves as (=) a "bath" for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called **thermalization.**

(183)However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.

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| (184)Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1). |
| (185)Bringing them to thermal equilibrium requires (e) an external bath. |
| (186) But so far, there has been no rigorous proof of (e) the existence of **such many-body localization**. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption. |
| (187)The concept of **localization** was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations). |
| (188)The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century. |
| (189)But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information. |
| (190)Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment. |
| (191)For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood. |
| (192) It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7]. |
| (193) Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature. |
| (194) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators. |
| (195)However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8]. |
| (196)Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems. |
| (197) Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL. |
| (198)As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked. |
| (199)His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice. |
| (200) An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site. |

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| (201)This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself). |
| (202)This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition. |
| (203)However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime: |
| (204) These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects. |
| (205)It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies). |
| (206) This is a realistic scenario for (e) all known systems. |
| (207)On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects. |
| (208)On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other. |
| (209)No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true. |
| (210) His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found. |
| (211)Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization? |
| (212)Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects? |
| (213)Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions. |
| (214)So in those systems, the quantum phase transition does indeed exist. |
| (215)But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension. |
| (216)Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters. |

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# PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION ATTRITION MODELS

## MODULE NINE

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Model is applied to following formulations Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society**

**Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice ($\alpha$=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP This paper is divided into four parts. In part A, some general considerations about (e&eb) gravitational radiation are followed by (e) a treatment of the scalar wave equation in the manner later to be applied to (e&eb) Einstein's field equations. In part B, a co-ordinate system is specified which is suitable for (e) investigation of outgoing gravitational waves from (e) an isolated axi-symmetric reflexion-symmetric system. The metric is expanded in (eb) negative powers of (e) a suitably defined radial co-ordinate r, and the vacuum field equations are investigated in detail. It is shown that the flow of information to infinity is controlled by (e&eb) a single function of (e) two variables called the news function. Together with initial conditions specified on a light cone, this function fully defines (eb) the behaviour of the system. No constraints of (e) any kind are encountered. In part C, the transformations leaving (eb) the metric in the chosen form are determined. An investigation of the corresponding transformations in (eb) Minkowski space suggests (eb) that no generality is lost by (e) assuming that the transformations, like the metric, may be expanded in (eb) negative powers of r. Gravitational Waves in General Relativity VII Waves from Axi-Symmetric Isolated Systems Royal Society Proceedings H. Bondi, M. G. J. van der Burg, A. W. K. Metzner Published 21 August 1962.DOI: 10.1098/rspa.1962.0161 STOP There is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant $\Lambda$, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive $\Lambda$ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the $\Lambda>0$ generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible. Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of (e) GW signals from BBHs with component spins $\chi_{1,2}\in[-0.99,0.99]$ aligned (e&eb) with the orbital angular momentum, component masses $m_{1,2}\in[2,48]M_\odot$, and (e&eb) total mass $M_{total}\leq 50\ M_\odot$. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin**

**bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins >+0.6 and component masses ≳20 M⊙. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201**

## 18. MODEL FORMULATION

Let $p_9(t, a_9)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(g_9)Q_9$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_9(0) = 0, \quad g_9(P_9) > 0, \ P_9 \ 0, \qquad \lim_{P_9 \to \infty} g_9(P_9) = (K_9)_9 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_9 + (\rho_a)_9 + D_9(a_9, P_9, Q_9)\rho_9 + 0, \qquad 0 < a_9 \ \infty,$$

$$\rho_9(t,0) = \int_0^\infty n_9(a_9, P_9, Q_9)\rho_9(t_9, a_9)da_9$$

Where $D_9$ is the mortality function and $n_9$ is the fertility function.

Put $_9$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_9}{dt}\right]_9 = x_9\left[\frac{dQ_9}{dt}\right]_9 - x_9\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_9}{dt}\right]_9 = x_3\left[\frac{dQ_3}{dt}\right]_9 - x_4\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_9}{dt}\right]_9 - x_6\left[\frac{dQ_3}{dt}\right]_9$$

Where $x_9, x_9, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_9$ and $g_9(P_9)Q_9$

Now when you add the term $(b_9Q_9)$ with $g_9(P_9)(Q_9)$ we are infect see is that $(b_9Q_9)$ is credited incrementally part of $g_9(P_9)(Q_9)$ or instant nearly, when subtracted from $\frac{dQ_9}{dt}$ , it means that $g_9(P_9)(Q_9)$ is slowly eaten by $\frac{dQ_9}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_9 = \mu_9(a_9) + (f_9)_9(P_9)_9 + (f_9)_9(Q)_9 \quad (2)$$

$$n_9 = \beta_9(a_9)h_9(Q_9), \quad (3)$$

Where $\mu_9$the age- dependent debit rate whit is $\mu_9(a_9)$ 0 $\forall a > 0$; $(f_9)_9$ is the density- dependent death rate with $f_9(0) = 0$, $(f_9^9)_9(P_9)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_9)_9(P_9) = (k_9)_9$ ; $f_9$ is the functional response of creditors to prey with $[f_9]_9(0) = 0, (f_9^9)_9(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_9)_9(Q_9) = [k_9]$ ; $\beta$ is the age dependent birth

rate with $\beta(a) \geq 0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h'(Q) \leq 0, \forall Q \geq 0$, and $\lim_{Q\to\infty} h(Q) = k_3 \geq 0$.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_9 + (\rho_a)_9 = -\mu_9(a_9) + (f_9)_9 P_9) + \big(f_9)_9(Q_9)\big)(\rho_9)_9 \tag{4}$$

$$\rho_9(t_9, 0) = h_9(Q_9)\int_0^\infty \beta_9(a_9)\rho_9(t, a_9)da, \tag{5}$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(P_9)Q_9 \tag{6}$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_9$ and $Q_9$ of (4-6) satisfy

$$0 \leq \rho_9(t, a_9) \leq \rho_9^*(t, a_9),$$

$$0 \leq Q_9(t) \leq Q_9^*(t),$$

Where $\rho_9^*(t, a_9)$ and $Q_9^*(t)$, are solutions of

$$(\rho_t^*)_9 + (\rho_a^*)_9 = -(\mu_9(a_9) + (f_9)_9(P_9^*) + (f_9)_9 Q_9^*))\,\rho_9^*, \tag{7}$$

$$\rho_9^*(t, 0) = \int_0^\infty \beta_9(a_9)\,\rho_9^*(t, a_9)da_9 \tag{8}$$

$$\frac{dQ_9^*}{dt} = -b_9 Q_9^* + g_9(P_9^*)Q_9^*, \tag{9}$$

With $P_9^*(t) = \int_0^\infty \rho_9^*\ (t, a_9)da_9$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_9^*(t)}{dt} = P_9^*(t)\big(r_9 - (f_9)(P_9^*) - (f_9)_9(Q^*)\big), \quad (10)$$

$$\frac{dQ_9^*(t)}{dt} = Q_9^*(t)(-b_9 + g_9(P_9^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_9(a)e^{-\int_0^a \mu_9(\tau)dt}e^{-ra}\ da_9 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_9^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(t)dt}\ da_9 > 1, \quad (H1)$$

Which guarantees $r_9 > 0$.

*Remark.* If $\beta_9(a) = [\beta_0]_9 a_9 e^{-\alpha_9 a_9}, \mu_9(a_9) \equiv \mu_9$ then the solution of (12) is

$$r_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_9]_9(P_9) = 0, [f_9](Q_9) = \lambda_9 Q_9\ h_9(Q_9) = 1, \quad and\ \ g_9(P_9) = c_9 P_9$$

$$[\rho_9]_9 + [\rho_a]_9 = -(\mu_9 + \lambda_9 Q_9)\rho_9$$

$$\rho_9(t,0) = \int_0^{\infty} [\beta_0]_9 a_9 e^{-\alpha_9 a_9} \rho_9(t,a_9)\, da_9$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + c_9 P_9 Q_9$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_9}{dt} = \omega_9 P_9 - \lambda_9 P_9 Q_9$$

$$\frac{dQ_9}{dt} = b_9 Q_9 - c_9 P_9 Q_9$$

Where $\omega_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$

THE EQUILIBRIUM $(\rho_0)_9(a_9), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_9(a_9), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_9 + (\rho_a)_9 = -\big(\mu_9(a_9) + (f_9)_9(P_9)\big)\rho,$$

$$\rho_9(t,0) = \int_0^{\infty} \beta_9\,(a_9)\rho_9(t,a_9)da_9 \qquad (13)$$

The solution of (13) has the form of

$$\rho_9(t,a_9) = Ce_9^{r(t-a_9)} - \int_0^a (\tau)d\tau - \int_0^9 (f_9)_9\big(P_9(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_9(t) = C_9 e^{r_9 t - \int_0^t (f_9)_9(P_9(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_9^{-9}(r)$ if $k_9 > r$. This is equivalent to

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{(k_9)_9 a_9} da_9 < 1.$$

In fact, if we define

$$R_9(s_9) = \int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau)d\tau} e^{-s_9 a_9}\, da_9$$

$R(s)$ is a decreasing function. $R_9(r_9) = 1$ And $R_9(k_9)_9 < 1$ imply $r_9 < (k_9)_9$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_9(a_9), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{k_9 a_9} da_9 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_9 \equiv \int_0^\infty (\rho_0)_{99}(a_9)da_9$ *satisfies*

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{-(f_9)_9(P_0)a_9} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_9(a_9), 0)$

Let $u_9(t,a_9) = \rho(t,a_9) - \rho_0(a_9)$ Then

$$N_9(t) \equiv \int_0^\infty u_9\,(t, a_9)da_9 = \int_0^\infty \rho_9(t, a_9)da_9 - \int_0^\infty (\rho_0)_9(a_9)da = P_9(t) - P_0)_9$$

And

$$(u_9)_9 + (u_a)_9 = (\rho_t)_9 + (\rho_a)_9 - (\rho_0^9)_9(a_9)$$

$$= -\mu_9(a_9) + (f_9)_9(P_9) + (f_9)_9(Q_9))\,(u_9 + (\rho_0)_9) + \mu_9(a_9)(\rho_0)_9 + (f_9)_9(P_0)_9(\rho_0)_9$$

$$= -\mu_9(a_9)u_9 - (f_9)_9(P_9) + (f_9)_9(Q_9))\,u_9 - (\rho_0)_9\,(f_9)_9(P_9) - (f_9)_9(P_0)_9$$

$$+ (f_9)_9(Q_9)), \qquad (15)$$

Where $(\rho_0^9)_9(a_9)$ denotes $d(\rho_0)_9(a_9)/da_9$

Expansion $(f_9)_9(P_9)$ and $(f_9)_9(Q_9)$ at $(P_0)_9$ and 0 yield respectively

$$(f_9)_9(P_9) = (f_9)_9((P_0)_9 + N_9)$$

$$= (f_9)_9(P_0)_9 + (f_9^9)_9(P_0)_9N_9 + o_9(N_9), \qquad (16)$$

$$(f_9)_9(Q_9)$$

$$= (f_9^9)_9(0)Q + o_9(Q_9) \qquad (17)$$

Write $f_9(P_0) = f_{90}, f_9'(P_0) = f_{90}'$, and $f_9'(0) = f_{90}'$ and substitute (16) and (17) into (15) to obtain

$$(u_9)_9 + (u_a)_9 \approx -(\mu_9(a_9) + (f_{90})_9)u - (f_{90}^9)_9(\rho_0)_9(N_9) - (f_{90}^9)_9(\rho_0)_9(Q_9)$$

$$u_9(t, 0) = \rho_9(t, 0)(\rho_0)_9(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t, a) + (\rho_0)_9(a_9)da_9 - (\rho_0)_9\,(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t, a_9)da_9 + (h_9(Q_9) - 1)(\rho_0)_9\,(0) \qquad (18)$$

Expanding

$$h_9(Q_9) = 1 + (h^9)_9(0)Q_9 + 0(Q_9) \equiv 1 + h_0'Q_9 + 0Q_9$$

And substituting it into (18), we arrive at

$$u_9(t,0) \approx \int_0^\infty \beta_9(a_9)u_9(t,a_9)da_9 + (h'_{a_9})(\rho_0)_9Q_9$$



PREDATOR-PREY POPULATION MODELS

The linearization is the following :

$$[u_9] + [u_9] = -(\mu_9(a_9) + [t_{90}]_9u_9 - [\alpha_9]_9(a_9)N_9 - [\alpha_9]_9(a_9)Q_9$$

$$u_9(t,0) = \int_0^\infty \beta_9(a_9)u_9(t,a_9)da_9 + [\xi_9]Q_9 \qquad (19)$$

$$\frac{dQ_9}{dt} = -b_9Q_9$$

where

$$[\alpha_9]_9(a_9) = [t_9^9a]_9[p,o]_9(a_9)d2(a_9) = [f_{90}^9]_9(\rho\, o\,)_9(a_9)\xi = [\mathrm{h\ o}^9]_9[\rho\mathrm{o}]_9(\mathrm{o}) \text{ and } \mathrm{b}_9 = \mathrm{b}_9 - g_9(P_9)$$

If $g_9(P_0) > b$, clearly, $(\rho_0)_9(a_9), 0)$ is unstable and the predator population is persistent. Assume $g_9(P_0)_9 < b$.

Then, $Q_9(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_9 + u_a = -(\mu_9(a_9) + (f_{90})_9)u_9 - (\alpha_9)_9(a_9)N_9$$

$$= -(\mu_9(a_9) + f_{90})u_9 - \alpha_9(a_9)\int_0^\infty u_9(t,a_9)da_9$$

$$u_9(t,0) = \int_0^\infty \beta_9\,(a_9)u_9(t,a_9)da_9$$

Using the method developed by Cushing [2, 3], we suppose

$$u_9(t,a_9) = \tilde{u}_9(a_9)^{e^{\gamma(t-a_9)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_9}\,da_9$. Then,

A solution of the first differential equation for $\tilde{u}(a_9)$ gives,

$$\tilde{u}_9(a_9) = e^{-\int_0^a \mu_9(\tau_9)d\tau-(f_{90})_9 a}\left(\tilde{u}_9(0) - w\int_0^a e^{\int_0^\tau \mu_9(s_9)ds+(f_{90})}\tau_9\,(\alpha_9)_9(\tau_9)e^{\gamma_9\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\right)\tilde{u}_9(0)$$

$$+w\int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\int_0^a (\alpha_9)_9(\tau_9)e^{-\gamma_9\tau_9}e^{(f_{90})_9\tau_9+\int_0^\tau \mu_9(s_9)ds}d\tau_9\,da_9$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}da\right)(\tilde{u}_9(0))_9$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\int_0^a e^{(f_{90})_9\tau_9+\int_0^\tau \mu_9(s_9)ds+y\tau}\alpha_9(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(t_9)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_9\,(a_9)Y_9(\gamma_9,a_9)da_9\right)u_9(0)$$

$$+w\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)\int_0^a(\alpha_9)_9(\tau)_9(Y^{-9})_9(\gamma_9,\tau_9)d\tau_9\,da_9=0,$$

(24)

$$\left(\int_0^{\infty}Y_9(\gamma_9,a_9)da_9\right)\tilde{u}_9(0)$$

$$+\left(-1-\int_0^{\infty}Y_9(\gamma_9,a_9)\int_0^a(\alpha_9)_9(\tau)_9(Y^{-9})_9(\gamma_9,\tau_9)d\tau_9\,da_9\right)w=0.$$

Equation (20) has a solution of the form of $u_9(t,a_9)=u_9(a_9)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_9(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_9(\gamma_9,a_9)=\int_0^a(\rho_0)_9(\tau_9)Y^{-9}\,(\gamma_9,a_9)d\tau_9$. We obtain the characteristic equation

$$f_{90_9}^{l}\left(\int_0^{\infty}Y_9(\gamma_9,a_9)da_9\right)\left(\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)I_9(\gamma_9,a_9)da_9\right)$$

$$+\left(1-\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)da\right)\left(1+(f_{90}^{l})_9\int_0^{\infty}Y_9(\gamma_9,a_9)I_9(\gamma_9,a_9)da_9\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_9(a_9) = (\rho_0)_9(0)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - (f_{90})_9 a_9}$$

$$I_9(\gamma_9, a_9) = \int_0^a (\rho_0)_9(\tau_9)e^{\gamma_9\tau_9 + (f_{90})_9\tau + \int_0^\tau \mu_9(s_9)ds_9}d\tau$$

$$= \frac{(\rho_0)_9(0)}{\gamma_9}(e^{\gamma_9 a_9} - 1)$$

Let

$$E_9(a_9) = e^{-(f_{90})_9 a_9 - \int_0^a \mu_9(\tau_9)d\tau_9}$$

Then $[Y_9]_9[\gamma_9]_9, a_9) = E_9(a_9)e^{-\gamma_9 a_9}$. by a simple calculation,

$$\int_0^\infty \beta_9(a_9)Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma}\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}\ da\right);$$

$$\int_0^\infty Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma_9}\left(\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}\ da_9\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_9(0)}{\gamma_9}[f_{90}^9]\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}da_9\ (\int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9)$$

$$+\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right).$$

$$\left(1+\frac{(\rho_0)_9(0)}{\gamma_9}(f_{90}^{'})_9\left(\int_0^{\infty}E_9(a_9)da_9-\int_0^{\infty}E_9(a_9)e^{-\gamma_9 a_9}da_9\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right)(1+\frac{(f_{90}^{'})_9}{\gamma_9}(\rho_0)_9(0)\int_0^{\infty}E_9(a_9)da_9)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (94), to

$$1-\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)da_9=1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\neq 0,$$

Since $\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)\int_0^{a}\alpha_9(\tau_9)Y^{-9}(\gamma_9,\tau_9)d\tau_9 da_9>0$. Therefore, the characteristic solution is

$$\gamma_9=-(f_{90}^{'})_9(\rho_0)_9\int_0^{\infty}E_9(a_9)da_9.$$

Since $f_9'(P_0)_9>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_9),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_9),0)$ which is related to the extinction of predators (debit summation types: that which acts as a resource)and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $( p^*(a), Q^* )$. It must satisfy the equations:

$$[p_a^*]_9 + (\mu_9(a_9) + [f_9]_9( P_9^*) + [f_9]_9(Q_9^*) )P_9^* = 0, \quad (26)$$

$$p_9^*(0) = h_9(Q_9^*) \int_0^\infty \beta_9(a_9)p_9^*(a_J)p_9^*(a_9)da_9 \quad (27)$$

$$g_9(P_9^*) - b_9 = 0, \quad (28)$$

Where $P_9^*=\int_0^\infty p_9^*(a_9)\, da_9$

Equation (28) yields $P^* = g^{-9}(b)$,

From (26),

$$p_9^*(a) = p_9^*(0)e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}$$

Substituting into (27), we need to solve

$$\left(1 - h_9(Q_9^*)\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}da_9\right)p_9^*(0) = 0$$

For $Q_9^*$

Fix $p_9^*(0)$ and $P_9^*$ Let

$$H_9(z_9) = 1 - h_9(z_9)\int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9}\, da_{,9}$$

And assume that

$$[k_3]_9 \int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(g_9^{-9}(b_9))a_9 - [k_9]_9 a_9}\, da_9 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_9(z_9)] > 0.$$

Clearly,

$$\frac{dH_9}{dz_9} = h_9(z)\left([f_9']_9(z_9)\right)\int_0^\infty a_9\beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - f_9(P_9^*)a_9 - f_9(z_9)a_9} da_{\square}$$

$-h'(z_9)\int_0^\infty \beta_9(a_9)\, e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9}\ da_9$

y the assumptions on $\left[f_9^9\right]_9 and\ [h^9]^9 dH_9/dz > 0$. Hence, if

$$H_9(0) = 1 - \int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9} da_9 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_9(0) < 0$ follows if $g_9^-(b_9) < [f_9^{-9}]_9(r_9)$, where $r_9$ is the solution of (12).

Indeed, since

$$R_9(s_9) = \int_0^{\infty} \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 e^{-y\frac{a}{9}} da_9}$$

Is a decreasing function of $s_9$ when $P_9^* = [g^{-9}]_9(b_9) < [f_9^{-9}]_9(r_9)$, that is $[f_9]_9(P_9^*) < r$, $R\big(f_9(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_9(g^{-9}(b))) > 1, i.e.,$

$\int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9\left(g_9^{-9}(b)\right)a_9} > 1.$ (H4)

Then the system (4)-(6) has a unique positive equilibrium $(P_9^*(a_9), Q_9^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_9^*(a_9), Q_9^*)$ is more complicated. We again consider perturbations $m_9(t_9 a_9) = p_9(t_9, a_9) - p_9^*(a_9)$ and $W_9(t) = Q_9(t) - Q^*$.Then

$M_9(t) \equiv \int_0^{\infty} m_9(t, a_9) da_9 = P_9(t) - P_9^*,$

$$[m_9]_9 + [m_a]_9 = [p_9]_9 + [p_a]_9 - \left[p^*\right]_9(a_9) \approx -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9$$

$$-[f_{9*}^9]_9 p_9^*(a_9)M - [f_{9*}^9]_9 p_9^*(\mathrm{a})W_9,$$

$$m_9(t,0)p_9(t,0) - p_9^*(0) \approx h_* \int_0^{\infty} \beta_9(a_9) m_9(t, a_9)\, da + \frac{[h_*^9]_9}{[h_*]_9}\, p_9^*(0)W_9$$

$$\frac{dW_9}{dt} = \frac{dQ_9}{dt} \approx [g_*^9]_9 Q_9^* M_{,9}$$

Where $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $[f'_{9*}]_9$ denoting $[f'_9(P^*)]_9 [f'_{9*}]_9$ denoting $[f_9]_9$ $(Q_9^*$, $h_9^*$ denoting $h'_*$ denoting $h'_9$ $(Q_9^*)$, and $[g'_*]_9$

Denoting $[g']_9(P_9^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_9]_9 + [m_a]_9 = -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9 - [f'_{9*}]_9 [p_9^*](a_9)M_9 - [f'_{9*}]_9 p_9^*(a_9)W_9$$

$$m_9(t,0) = [h_*]_9 \int_0^\infty \beta_9(a_9)m_9(t,a_9)da_9 + \frac{[h_*^9]_9}{[h^*]_9}[p^*]_9(0)\,W_{,9} \qquad (30)$$

$$\frac{dW_9}{dt} = [g_*^9]_9[Q_9^*]M_9$$

Integrating $[m_9]_9(t,a_9)$ with respect to a from 0 to ∞, we have

$$\frac{dM_9(t)}{dt} = \int_0^\infty [mt]_9(t,a_9)da_9$$

$$= -\int_0^\infty [m_a]_9\,(t,a_9)da_9 - \int_0^\infty \mu(a_9)m(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + [f'_{9*}]_9 \int_0^\infty p_9^*(a_9)da_9\right)M_9 - [f_{9*}^9]_9 \int_0^\infty p_9^*(a_9)da_9 W_9$$

$$= m_9(t,0) - \int_0^\infty \mu_9(a_9)\, m_9(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + [f'_{9*}]_9 P_9^*\right)M_9 - [f'_{9*}]_9 P_9^* W_{.9}$$

From (30), we arrive at the equations

$$\frac{dW_9(t)}{dt} = \int_0^{\infty} ([h_*]_9\beta_9(a_9) - \mu_9(a_9))m_9(t,a_9)da_9$$

$$-\left([f_{9*}] + [f_{9*}]_9 + \left[f_{9*}^{'}\right]_9 P_9^*\right) M_9 + \left(\frac{\left[h_*^{'}\right]_9}{[h_*]^9}\ P_9^*(0) - f_{9*}^{'} P_9^*\right) W_9,$$

(31)

$$\frac{dW_9(t)}{dt}\left[g^{'}{}_*\right]_9 Q_9^*\, M_{.9}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{"*}]_9\beta_9(a_9) - \mu_9(a_9))m_9(t,a_9)da_9 = ([h_*]_9\beta_9(\theta_9) - \mu_9(\theta_9))M_{,9}$$

Where $0< \ \theta_9 = \theta_9(t) < \infty$.

And if $a_9^9[a_9^9]2<4\ [a_9]2<4\ a_9e$,

$\text{Re}\,\lambda_9[\lambda_9]_9 = \text{Re}\lambda_9[\lambda_9]_9 = -\frac{9}{9}a_9[a2\ ]_9 < 0.$

Let $\Delta=\max_{9<u<\infty}\left|h_*[h_*]_9\beta_9(a_9) - \mu_9(a_9)\right|$. Then$|n_9(t)|$ ⍰⊿∀t ⍰0. Hence we obtain the following result.

**THEOREM 4**

Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_9^9 <4\ a_9e$ and $[a_9] > 1\Delta$,or $[a_9^9]$ ⍰4$[a_9]e$ and $(a_9 - \sqrt{(a_9^9)_9 - 4(a_9)_9e) > 1\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+9}\right)$ and $\beta_9(a_9) = ae^{-ma}$ Then $\Delta=[\mu_0] + 2$ Simple calculation leads to

$$\int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9) + d\tau_9 da} = \frac{1}{(m + \mu_0 + 1)^9} + \frac{1}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_9 + [\mu_0]$ 0.5. Assume$[f_9](x_9)= K_9 x_9$ Then the hypothesis (H9) is satisfied We take the conversion rate function $g_9(x_9)= [0x_9]/(2 + x_9)$; the functional responses $[[f_9](x) = [0]\, x_9/(1 + x_9)$and $h_9(x_9)=(1+\omega_9 x_9)/(1 + x_9)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_9 + [\mu_0] + [K][P_0])}{[m_9]_9 + (\mu_0)_9 + [K][P_0])_9^3\ [K]_9[P_0]_9^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_9$i e $[P_0] < b/(\delta_9 - b_9)$then the equilibrium$(P_0)$ $[P_0]_9(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_9$

$[f_{9*}] = K_9\ P^*\ [f_{9*}] = \frac{\sigma Q_9^*}{9+Q_9^*}$ $[h_*] = \frac{9+\omega_9 Q_9^*}{9+Q_9^*}$

$[f_{9*}] = K_9\ [f_{9*}^9] = \frac{\sigma_9}{(9+Q_9^*)^9}\ [g_*^9] = \frac{\sigma_9}{(9+P_9^*)^9}$ $[h_*^9] = \frac{\omega_9 - 9}{(9+Q_9^*)^9}$

In order to have (H3) and (H4) hold,$m, \mu_0, K, b, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_9}{(m+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9+b_9}+\sigma)^9}+\frac{1\omega_9}{([m_9]_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9+b_9}+\sigma_9)^3}<1$$

And

$$\frac{1}{(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9})^9}+\frac{1\omega_9}{(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9}+(m_9+[\mu_0]_9+\frac{K_9 b_9}{\delta_9-b_9}+\sigma_9)\,Q_9^*)^3$$

$$=m_9+[\mu_0]_9+3+\frac{K_9 b_9}{\delta_9-b_9}(1+Q^*)(1+\omega_9 Q_9^*$$

Now

$$\begin{bmatrix}1\\1\end{bmatrix}-4\begin{bmatrix}a_9\end{bmatrix}e=\left(\frac{1K_9 b_9}{\delta_9-b_9}\right)^9+\frac{4K_9 b_9\sigma_9\,Q_9^*}{(\delta_9-b_9)(9+Q_9^*)}+\frac{\sigma_9^9 Q_9^{*9}}{(1+Q_9^*)^9}$$

$$\frac{4\,b_9\sigma_9\,(\delta_9-b)Q^*}{\delta_9(1+Q_9^*)^9}-\frac{4\,(\delta_9-b_9)^9(1-\omega_9)Q_9^*}{\delta_9(1+\omega_9 Q_9^*)(1+Q_9^*)}P^*(0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_9 \, ae^{mu}_9 \quad [\mu_0]_9 - \frac{a_9}{a_9+9}$$

Hence if $[a^9_9 > 4(a_9)_9 e$ and

$$[a_9] = \frac{1K_9 b_9}{\delta_9 - b_9} + \frac{\sigma_9 Q^*_9}{1 + Q^*_9} > 1\Delta$$

Some of the examples to which the model has been applied are:

**New look at the galaxy-shaping power of black holes**

(120) "It creates (eb) these bubbles, and the bubbles keep the plasma hot.

(121) That's what prevents galaxies from (e) becoming even bigger than they are now."

(122) Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.

(123) In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.

(124) Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.

(125) "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.

(126) "It's like nature's thermostat, which keeps these galaxies from (e) growing.

(127) If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."

**A new look at the galaxy-shaping power of black holes**

(128) An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the

(129) McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.

(130) Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.

**Doomed satellite missions**

(131) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.

(132) At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.

(133) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.

(134) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.

(135) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.

(136) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people

(137) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."

(138) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.

(139) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(140) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.

(141) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.

(142) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.

(143) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.

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| (144) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well. |
| (145) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma. |
| (146) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6. |
| (147) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus. |
| (148) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (149) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (150) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (151) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (152) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (153) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (154) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (155) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (156) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (157) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (158) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (159) We know dark matter exists only from (e) its gravitational effects. |
| (160) It holds together galaxies and (e&eb) galaxy clusters. |
| (161) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (162) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |

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| (163) | They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (164) | X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (165) | They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (166) | This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (167) | An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |
| (168) | Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (169) | "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (170) | The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (171) | Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (172) | Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (173) | "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (174) | Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (175) | "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (176) | Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) |

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| consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories. |
| (177) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen. |
| (178) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer. |
| (179) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants. |
| (180) These were treated as (=) "aware" and "unaware" trials, respectively. |
| (181) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it. |
| (182) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another. |
| (183) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion. |
| (184) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness. |
| (185) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said. |
| (186) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness." |
| (187) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner. |
| (188) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain. |
| (189) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said. |
| (190) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

(191) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter.

(192) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined.

(193) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions.

(194) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.

(195) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.

(196) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.

(197) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.

(198) Noda and his colleagues instead studied the change in emissions over (e&eb) time.

(199) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.

(200) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).

(201) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.

(202) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.

(203) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson’s problem**

(204) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking “palm tree” organization.

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| (205) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society |
| (206) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible? |
| (207) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem. |
| (208) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other. |
| (209) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible. |
| (210) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges. |
| (211) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century. |
| (212) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching. |
| (213) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom. |
| (214) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth. |

**Spheres on trees**

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| (215) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration. |
| (216) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement. |
| (217) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum. |
| (218) By visualizing the problem in this way, the researchers noticed that these particular graphs don’t (e) have lots of branches extending from lots of (e) other branches. |
| (219) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure. |

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| (220) The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network. |
| (221) As a result, it provides (eb) an important clue for finding the global minimum. |
| (222) It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum. |
| (223) This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large. |

**Thompson's problem 2**

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| (224) By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society |
| (225) The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima. |
| **(226)** The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node. |

**Implications for Thomson's and Smale's problems**

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| (227) Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before. |
| (228) "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around." |
| (229) As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem. |
| (230) "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view." |
| (231) In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges. |
| (232) From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear. |
| (233) "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK. |
| (234) "For larger numbers, there will likely be (=) multiple funnels. |
| (235) Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum." |

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| (236) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem. |
| (237) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said. |
| (238) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org |

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| The nontrivial equilibrium is stable. |
| IF we define $[a_9]_9 = [f_{9*}]_9 + [f_{9*}]_9 + [f_{9*}]\,P^*_9$ |
| $$[a_9]_9 = \left[f^{9}_{9*}\right]P^*_9 - \frac{h^9_*}{h_*}P^*(0)$$ |
| $$e=[g^9_*]_9\,Q^*_9$$ |
| $$\eta_9(t) = \left[h_*\right]_9\beta_9\left(\theta_9(t)\right) - \mu_9(\theta(t)),$$ |
| The system(32)can be written as the following system: |
| $$[M(t)]\left[a_9\right]M_9 - [a_9]_9W_9 - [n(t)]M_9$$ |
| $$w_9(t) = eM_9$$ |
| Write |
| $$A=\begin{pmatrix}-[a_9]_9 - [a_9]_9 \\ e \qquad 0\end{pmatrix}$$ |

$$B=(t)=\begin{pmatrix} [\eta(t)]_9 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_9 \\ W_9 \end{pmatrix}$$

The solution of (32)is

$$X_9(t) = e^{A,t}[X]_9(0) + \int_0 e^{At}$$

**Note: There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and**

**sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by**

**term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation**

**alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation,**

**appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module**

**LEN NUNNEYS**

## Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

## Module numbered one

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems

and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \quad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1965); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be referred to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 6985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_6}{dt} = [N(t)]_6[(F_6)(R(t)]_6 - [M]_6[R(t)]_6$$

$$\frac{d[R(t)]_6}{dt} = [B]_6\,[(R)_6(t-T_6)] - [D]_6[(R(t)]_6 - [N(t)]_6\left[G\left(R(t)\right)_6\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_6}{dt} = a_6[N(t)]_6 - a_6[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_6 = a_3[N(t)]_3 - a_4[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_6 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

irst, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.6b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The

regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 6a is at the

**TIME DELAY**

FIG. 6. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.6$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**DELAY DUE TO PREDATOR MATURATION**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_6}{dt} = N_6(t - T_6)F_6(R_6(t - T_6) - N_6(t)M_6((R_6(t) \qquad (5)$$

$$\frac{dR_6(t)}{dt} = B_6\big(R_6(t) - D_6(R_6(t)\big) - N_6(t)G_6(R_6(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_6}{dR_6} + \frac{N_6 dG_6}{dR_6} - \frac{dB_6}{dR_6} > \frac{N_6 G_6}{2M_6}\left[\frac{dF_6}{dR_6} + \frac{dM_6}{dR_6}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_6/dR_6$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_6$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_6$) and the other derived from the non-delayed or "immediate" processes($I_6$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

ABSOLUTE STABILITY

$$L_6 = \frac{N_6 G_6}{2M_6}\frac{d(6nF_6)}{dR_6} (\geq 0) \qquad (6a)$$

$$I_6 = -\frac{N_6 G_6}{2M_6}\frac{d(6nM_6}{dR_6} (\geq 0) \qquad (6b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_6 > L_6 - I_6$, must be modified ahe $dM_6/dR_6$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_6 > L_6 - I_6, \quad \text{when} \quad P_6^2/4 > I \qquad (8a)$$

$$P_6 > L_6 - I_6 + \frac{(I_6 - P_6^2}{(L_6 - I)} \quad \text{when} \quad P_6^2/4 < I. \quad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

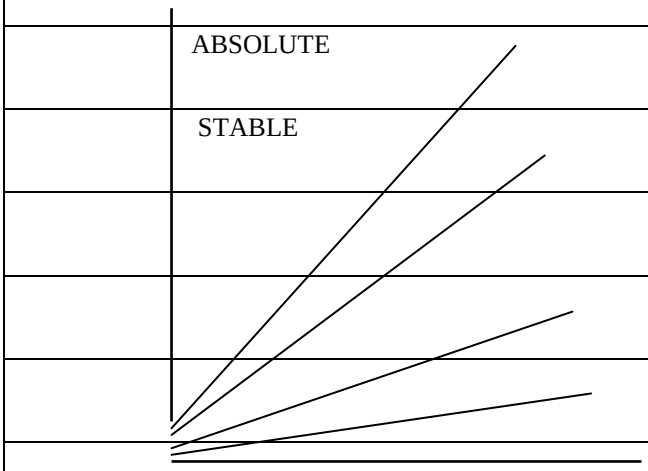



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

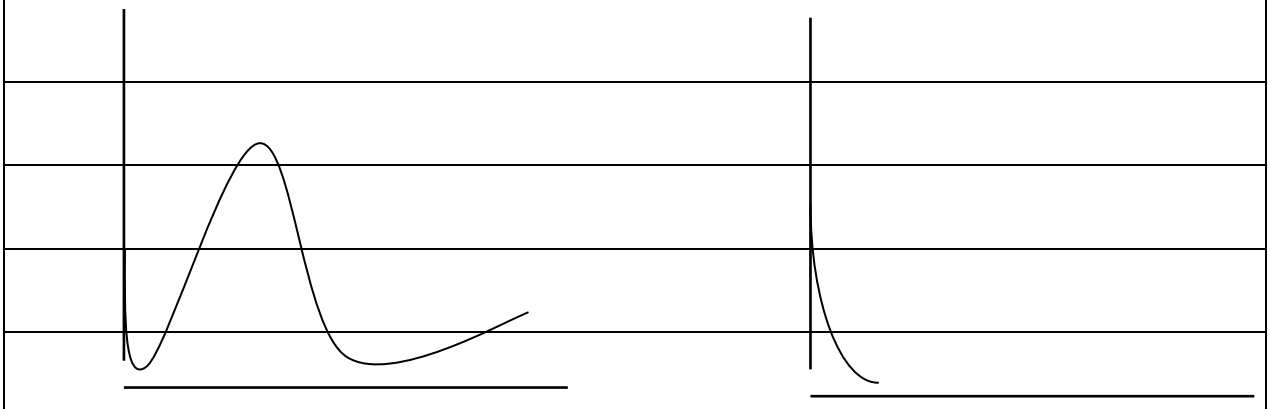
Can never completely offset the lagged effect,$L_6$; condition (8b) ensures that the total debit regulation, $P_6$, must always be positive even when $I_6$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_6$ and $L_6$. As $I_6$ increases relative to $L_6$ the zone of stability is enlarged; however,even when $I_6$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_6$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



**TIME DELAY**

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (6). Parameters in each case M=6; $dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6 = 0.05$ For (a) $NGdF/dR = 0.3, -N_6G_6dM_6/dR_6 = 0.6$; for (b)$NG_6dM_6/DR_6 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_6t}{dt} = N_6(t - T_6)F_6\big(R_6(t - T_6)\big) - N_6(t)M_6\big(R_6(t)\big)$$

(9)

$$\frac{dR_6t}{dt} = B_6\big(R_6(t - T_6)\big) - D_6\big(R_6(t)\big) - N_6(t)G_6\big(R_6(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) and the effective prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we

could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

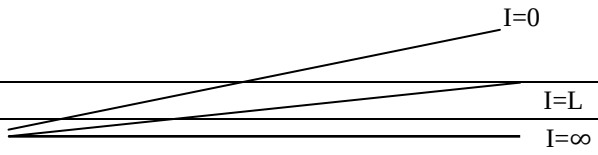
Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_6/dR_6$ is a measure of the control acting through the debit delay and setting $dB_6/dR_6 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_6/dR_6$ is a measure of the control acting through the credit delay although setting $dF_6/dR_6 = 0$ does not negate the effect of a predator delay; the value of $M_6$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_6/dR_6$and normalized by $M_6$ (see6a). The absolute stability of such a system is summarized in Fig. 6 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

TOTAL

ABSOLUTELY STABLE

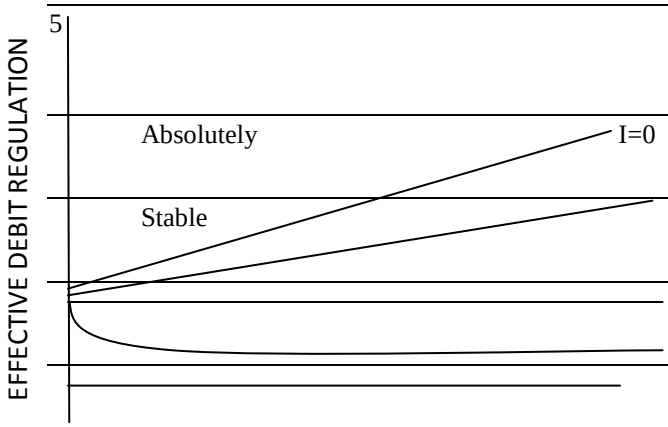



DELAYED DEBIT EFFECT

FIG6. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_6/dR_6 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.



-50 10

DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_6/dR_6 + N_6 dG_6/dR_6 + dB_6/dR_6)/M_6$). The immediate ($I_6$) and lagged ($L_6$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_6$(see 6b), is increased relative to the lagged part, $L_6$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_6 > I_6$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_6$is somewhat large than $I_6$ then the condition constraining prey regulation of recruitment ($dB_6/dR_6 < 0$) is

$$\frac{dD_6}{dR_6} + N_6\frac{dG_6}{dR_6} + \frac{dB_6}{dR_6} > \frac{(L_6 - I_6)M_6}{2} \qquad (10)$$

If $L_6 < I_6$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_6$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_6$

The - $D_6$ positions are defined by letting $\lambda_6 = iZ_6$ and separating (A1.2) into real and imaginary parts:

$$Z_6^2 + V_6 Z_6 \int_0^{\infty} \sin(Z_6 T_6) f(T_6) dT_6 - (X_6 + Y_6) = 0$$

$$Z_6 \left( U_6 - V_6 \int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6 \right) = 0$$

Noting that $X_6 + Y_6 = 0$ is we assume X+Y=0.It follows that $Z_6 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_6 = V_6 \int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6$$

If $U_6 > |V_6|$then given the constraints of $f(T_6)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A6.5)remains

greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

## ABSOLUTE STABILITY WITH A DELAY DUE TO CREDIT CONSOLIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_6^2 + (U_6 - V_6 + M_6)\lambda_6 - M_6\lambda_6 e^{-T_6\lambda_6} + [(U_6 - V_6)M_6 + Y_6] - [((U_6 - V_6 M_6 -]e^{-T_6\lambda_6} = 0$$

(For definitions see Appendix6) We follow the technique outlined in Appendix 6 and let $\lambda_6 = iZ_6$ and separate (A2.6) into real and imaginary parts

$$Z_6^2 + M_6 Z_6 \sin(Z_6 T_6) - (A_6 M_6 + Y_6) + (A_6 M_6 + X_6)\cos(Z_6 T_6) = 0$$

$$(A_6 M_6) Z_6 - M_6 Z_6 \cos(Z_6 T_6) + (A_6 M_6 - X_6)\sin(Z_6 T_6) = 0$$

Where $A_6 = U_6 - V_6$Using (A2.3), substitute for $Z_6^2$ and $Z_6$in (A2.2):

$$W_6^2(X_6^2 - 3M_6 A_6 X_6 - M_6^2(X_6 - Y_6) + 2M_6^3 A_6 + 2A_6^2 M_6^2]$$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_6 = 1 - \cos(Z_6 T_6)$.

In the absence of a delay the system is stable for $A_6 > 0,$ and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_6 > (X_6 - Y_6)/2\ M_6$ Thus a necessary condition for absolute stability is

$$A_6 > [\ X_6 - Y_6/2\ M_6$$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_6$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_6$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_6, A_6) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \leq W_6 \leq 2$ *is* of interest First consider the extreme values

$$Q_6(0, A_6) = A_6^2\ (X_6 - Y_6)$$

$$Q_6(2, A_6) = (A_6 + 2\ M_6)^2\ (2A_6 M_6 - (X_6 - Y_6))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_6}{\partial A_6} = W_6^2\left(2M_6^3 + 4A_6 M_6^2 - 3M_6 X_6\right) + W_6\left(3M_6 A_6^2 - 2A_6 X_6 + 2M_6\left(2X_6 + Y_6\right)\right)$$

$$+2A_6\left(X_6+Y_6\right)$$

$=2M_6^2W_6^2(M_6+A_6)+M_6A_6^2W_6+\left(2A_6M_6-X_6+Y_6\right)(A_6W_6+M_6W_6^2$

$+(2-W_6)((2M_6X_6+M_6Y_6)W_6+A_6(X_6+Y_6)),$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_6$quadratic in $W_6$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_6$has no minimum in the range $0 \le W_6 \le 2$.Defining

$Q_6\left(W_6,A_6\right)=a_6W_6^2+c_6 where\ a_6\ b_6\ and\ c_6$ Are functions of $A_6$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_6 \ge b_6 \ge 0$$

And

$$b_6^2 \le 4a_6c_6$$

$$Case(i): Y_6=0$$

$Let\ Y_6 = 0$ And try the substitution A= $X_6 = /2M_6\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_6 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_6 X_6 /$ $2M_6$) is positive for *W<2. Since* $Q_6$ increases with $A_6$(see A2.9), increasing $A_6$will generate no further solutions of $Q_6$=0. Hence the condition A>X/2M defines absolute stability when $Y_6 = 0$.

$Case\ (ii)$: $Y_6 = y_6\ X_6$ ,

$Let\ Y_6 = y_6\ X_6$And try the substitution $A_6 = X_6/(2 + A_6\ y_6)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_6/(2 + n_6 y_6)$ we have

$$W_6^2\ y_6\left[X_6 n_6 \frac{(1 + n_6 y_6)}{(2 + n_6 y_6)} + M_6^2\left(2 - n_6 + n_6 y_6\right)\right]$$

$$-W_6\ X_6\left[\frac{X_6(1 + n_6 y_6)}{M_6^2(2 + n_6 y_6)^2} + 2\ y_6\left(n_6 - 1\right) + \frac{X_6^2(6 + y_6)}{M_6^2(2 + n_6 y_6)} = 0.\ (A2.12)\right.$$

Since $n_6$ and $A_6$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_6$

Defining $E_6 = (2 + n_6 y_6), the\ condition (A2.66?) can\ be\ written\ as$

$[\frac{X_6(E_6 - 6)}{E_6^2}]^2[\frac{X_6^2}{M_6^2} - 4\ y_6\ E_6^2\frac{X_6}{M_6^2} + 4y_6(n_6 - 2)E_6^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in ($X_6 / W_6^2$) can be expressed as

$2E_6[y_6E_6+(y_6E_6(6+y_6)(E_6-n_6))^{6/2}]>\frac{X_6}{M_6^2}>2E_6[y_6E_6-y_6E_6(6+y_6)(E_6-n_6))^{6/2}]$

**ABSOLUTE STABILITY**

Hence if $X_6/M_6^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0$W_6$2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_6-6)}{E_6^2}\left(\frac{X_6}{M_6^2}\right)^2-2y_6\frac{(E_6+n_6^2y_6)}{E_6}\left(\frac{X_6}{M_6^2}\right)-4y_6(E_6-n_6)$$0

And has the solution

$\frac{X_6}{M_6^2}>\frac{E_6}{(E_6-6)}[y_6(E_6+n_6^2y_6)+[y_6^2(E_6+n_6^2y_6)^2+4y_6(E_6-6)(E_6-n_6)]^{6/2}]$

Given $E_6>n_6$this lower limit to $X_6/M_6^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_6/M_6^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_6$in the appropriate range and this limit is further lowered by (A2.64)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_6<n_6(ie.,2A_6M_6^2<X_6-Y_6$)then (A2.6)and

(A2.8)show that $Q_6$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_6}{M_6^2} > 2E_6[y_6E_6 - (y_6E_6(6 + y_6)(E_6 - n_6)]^{6/2}]$$

This equations throws light on the individual general ledger and Nature's general ledger as we call it : see various papers on the issue. powers, which are the energies of Matrka we have already discussed, are not the only obstacles the yogi must overcome. He must, for example, also resist the temptation to rest content with the miraculous yogic powers (siddhi) he acquires in the course of his spiritual develop- ment. Again to do this he must practice Yoga. Similarly, in order to pervade the Fourth State gradually through the other states in the manner proper to practice at the Individual level, the method is the same. He must practise the higher yoga of the Tantras which, turning his mind inwards and freeing it from discursive representations, allows him to penetrate into the Supreme Principle. 312 Once the yogi has attained this contemplative state, his main problem is to make it permanent. In the introverted state the gross external movement of the breath is suspended and with it the activity of the intellect, mind, individualised consciousness, powers of the senses and the ego. 313 When the yogi rises out of this state, he is liable to fall again into the lower order of creation generated by Maya if he does not maintain his awareness of the higher reality he has experienced and allows his awakened, illumined insight to be obscured by the dream-like vision of thought-constructs. 314 Naturally, the yogi must rise out of the introverted condition of suspension. It is inherent in the very nature of reality that it should move out of itself. 315 Pure, universal consciousness initially transforms itself into the vital breath 316 charged with the impression iyasana) of the power of awareness attained through introversion. By attending to the pulse (spanda) of the breath as it moves out of the absolute, the

yogi can develop an intuitive sense of the inherent unity of all he will perceive in the mental and physical spheres created by the outpouring of consciousness. In this way he realises that his own nature is everywhere present in all he perceives and that all things thus reside within him. Blessed with this insight his consciousness remains free and unlimited even at the individual level where the breath, mind, senses and body are active. If the yogi fails to do this, he finds himself once again beset by the strictures of his embodied existence and must, as before, try to pervade all his other states of consciousness with the aesthetic delight (rasa) and wonder (camatkara) of the Fourth State he experienced in contemplation. Again this means that he must strengthen his pure, empowered awareness that his universal nature manifests as all things. 317 In this way he discovers Siva's presence in every sphere of individualised consciousness ranging from the breath to outer objectivity. The yogi's mind then becomes tranquil and undistracted because wherever it may wander, the yogi perceives only Siva, his authentic nature. 318 Consciousness is thus freed of all external referents and the yogi's subjectivity is purified of all identification with the body or anything else that belongs to the objective sphere. The yogi then becomes detached from the opposites of pleasure and pain and is transcendentally free (kevalin). 319 yogis are again, however, liable to fall if he allows himself to get entangled in the play of opposites. This fall is more serious than the others because, although he is caught by the confining restrictions of individual- ised consciousness as before, he is now also affected by karma. Fleeing from pain in the pursuit of pleasure he is bound to act (karma) to minimise one and maximise the other and so is thrown down to the lowest level of embodied subjectivity (sakala). In order to regain his lost state, he must ascend gradually, by Siva's grace, from one order of subjectivity to the next and so free himself progressively of the limitations of the lower levels to gain the greater freedom and expansion of the higher. As he progresses, the objective sphere also evolves from the grossest perceptions of physical objects outside the lowest order to subjectivity, through to the subtler inner, mental

perceptions to finally reach the order of subjectivity that contains objectivity within itself and is free to externalise it at will. 320 The degree to which this process develops depends, as before, on the yogi's awareness of the Fourth State. In consonance with the general principle that the remedy should suit the defect, the yogi is instructed to seek this higher state of consciousness in the wonder (camatkara) or delight (ananda) he feels in moments of intense physical pleasure. At first he experiences this subtle consciousness for an instant in the subjective sphere. If he manages to catch hold of it, it becomes more intense as the cognitive and objective spheres are also gradually pervaded and vitalised by it. Occasions for this practice are, for example, the sense of satisfaction one feels after a good meal or the aesthetic delight one experiences when listening to good music or the pleasure of sexual union with the Tantric consort or even solitary sexual excitation. In these moments of delight the yogi can penetrate momentarily into his own authentic Siva-nature (sambhavavesa) through the empowered contact (saktasparsa) m he makes with it in the freedom of the pure subjectivity of the Fourth State. 322 If the yogi develops his awareness of this higher level of consciousness and maintains it, he eventually experiences it constantly. 323 Clearly, what prevents the yogi from attending to his state of consciousness rather than the circumstances which induce it is the craving for pleasure (abhila§a) born of ignorance — the source of every impurity which clouds consciousness. Craving directs the yogi's attention towards outer, worldly things and so he is caught in the net of thought- constructs. 324 To free himself of his worldly desires and reverse this binding extroversion, the yogi must eradicate its cause. To be freed of all the ups and downs of the path and no longer be tormented by the possibility of a fall, the yogi must see reality perfectly and completely. This insight is itself liberation and the moment it dawns the yogi is instantly freed. This sudden realisation is the goal of Tantric Yoga. E Tantra declares: "He who perceives reality directly, even for the brief moment it takes to blink, is liberated that very instant and never reborn again." 325 Although the yogi's body and mind

continue to function as before, they are like mere outer coverings 326 which contain, but do not obscure, the mighty, universal consciousness which operates through them. The yogi's body is (=) the universe, the senses (=) the energies that (eb, eb+)vitalise it, his mind (=) Mantra, the rhythm of his breath (eb) the pulse of time and his inner nature pure, dynamic consciousness. Raised above all practice, and hence all possibility of falling to lower levels, the yogi realises that he has always been free 327 and that his journey through (e&eb) the **dark land of Maya was nothing but a dream, a construct of his own imagination** The channels of the senses by which the yogi enjoys the pleasure of the objects of sense are those through which he fills the three worlds with the consciousness emitted from his own Heart. 78 **The Doctrine of Vibration T**he senses are outer, subsidiary cycles of conscious energy (anucakra), revolving around the core or primary cycle (mukhyacakra) of pure consciousness. They scintillate with the diversity of sensations and resonate in the field of their awareness like the feathers of a peacock unfurled and folded back in the ecstasy of its dance. 79 Ultimately it is Siva Himself Who, for His own pleasure, 80 extends and withdraws them, 81 not the individual soul. The outward movement of the senses is the creative flow of consciousness from the subject who, like an ocean, is the source of the waves of their activity and the final resting place of the flow of the rivers of sensations and thought-forms that unfold through the arising of the power of each of the senses. 82 The Stanzas on Vibration declare: Indeed the individual soul does not activate the impulse of the will [which directs the body's activity] by himself alone but through his contact with [his] own [inner] strength {bald) made in such a way that he identifies with it [thus acquiring its power]. 83 The close contact between the senses and universal consciousness vitalises them and sets them in motion. They share in the properties of consciousness, just as a heated iron ball can burn, cook or give out light. 84 Spanda doctrines does not agree with the Sarnkhya view that the individual soul (purusa) alone stimulates activity in the objective sphere (prakrti) by the contact he makes with it. The soul can do

this only because universal consciousness animates him. Moreover, consciousness, which has the power to impel this activity, cannot itself be inactive as the Sarnkhya maintains. The influence it has on lifeless matter implies an active exertion. 85 This activity is the subtle, inner outpouring of the Self (svatmocchalatta) within the Self. Outside the confines of time and space, in the centre of the circle of the senses, this inner activity initiates the expansion (prasara) of the flow of conscious energy. As the pure exertion (udyoga, udyama) of Being (bhava), it is the source of the initial unfolding (unmesa) of Becoming. 86 Ksemaraja explains: [By] 'exertion' (udyama) is [meant] the emergence of the supreme [level of] intuition (pratibha) which is the sudden outpouring of the [immanent and transcendental] unfolding of the reflective awareness of pure consciousness. 87 The direct, intuition (pratibha) of reality, free of thought-constructs, is the supreme power (parasakti) of consciousness. The yogi can gain Divine **Body and Sacred Circle of the Senses 153** this pure intuitive awareness by grasping the initial expansion (unmesa) of consciousness which engenders the emergence of thought and the activity of the senses. It can be discovered in the subtle Centre, between one thought and the next. The Stanzas on Vibration teach: The expansion of consciousness that takes place when one is engaged in a single thought should be known to be the source from whence another arises. One should experience that for oneself! 88 This does not mean that the unfolding expansion of consciousness {unmesa) is an intermittent, transitory occurrence. 89 As the source of all thought, including the notions of past, present and future, 90 it stands outside the confines of time. Consciousness is in a state of perpetual expansion. It is an ever developing process, a tiny part of which we observe as the flow of events around us and the thoughts and feelings within us. The totality of this process can never be grasped objectively. Objectivity is a state of fragmentation (bheda) in which discrete elements are discernible and mutually distinguished as individual, specifically definable entities. It is the realm within which thought and language operate. To grasp reality in its completeness, we must go

beyond the partial representations of thought and speech. To experience the primordial source and basis of all things, we must pierce through the outer periphery of thought and plunge into the Centre, to discover the instant in which thought, and with it, the sense of diversity, initially emerges. Ksemaraja cites a practice from the Vijnanabhairava through which the yogi can achieve this state of awareness: Checking [the movement of] attention (cit) once it has quit its object, it must not [be allowed] to move to any other; contemplation (bhavana) then blossoms forth by [experiencing] the state in the centre between them (tanmadhyabhava). 91 Left unattended and undisciplined, awareness spontaneously shifts from one object of thought to (e&eb) the next. The transition from one to the other entails a movement through (e&eb) a state of pure indeterminate (nirvikalpa) awareness divested of all thought-forms and perceptions. The yogi who succeeds in checking the movement of his attention, experiences within himself the subsidence of thought into the emptiness of the Centre which unfolds full of the powerful pulsation of consciousness within which all objectivity merges and becomes one 92 in the trans- cendental outpouring (lokottarollasa) of its aesthetic rapture. 93 turning in on himself, the yogi must exert his power of awareness to (e) discern the vitality of (e) his own blissful Spanda nature. Just as a distant object becomes (=) more clearly visible by (e) exerting an effort to (e) perceive it, in the same way, through (e&eb) practice, the yogi's own Spanda nature becomes (=) progressively more evident to him, 94 The effort the yogi exerts to free himself of ignorance is essentially one with the force of awareness that impels the senses and mind. This inner exertion is the basis of every means to liberation. Ultimately bestowing the bliss of self- realisation, it carries the yogi along the path to salvation with all its twists and turns like wind which blows a cloth in various directions. 95 As the basis and ultimate goal of all yogic practice, this exertion is the pure wonder (camatkdra) inspired by the realisation of the essential identity of consciousness (cit) and being (sat). 96 As such it is the divine radiance (sphuratta) of the Heart of one's own consciousness 97 where the powers of

will, knowledge and action merge in the harmony (samarasya) of bliss. 98 Thus the highest form of worship (varivasya) is the contemplation of the inner strength (bald) or universal exertion of one's own true nature. 99 "One should examine [that reality]," says Ksemaraja, "with the exertion which is Bhairava's nature and is true service (sevana) to the integral inner nature [of all things] and the unfolding of the activity of one's own vitality (ojas) which withdraws all duality [back into consciousness]." 100 Mindful of the eternal joy (ahlada) of his inner strength, the yogi must give up all desire for the fleeting moments of petty pleasure (ksanikasukhalava) ]0] he may glean from outer objects. He must even give up the desire to achieve enlightenment and a reality (tattva) which can never be realised through a personal act of will. 102 In short, he must die to himself in an act of profound faith (sraddha) and adoration (adara) to find rest in the supreme reality free from the binding dichotomy of means and objectives (upeyopayabhava) — the reality of his perfected effort. The yogi, now full of devotion for this reality unfolding within himself and the universe (antarmukhatattva), merges into it and, by achieving the pure consciousness and bliss (cidananda) of the power (bala) of his own nature, achieves oneness. 103 The yogi is firmly established in his true nature by recognising the creative unfolding (unmesa) of consciousness through which he is able to know all things, whether in the present, past or future. 104 Master of the vitality which gives life to his mind and body, he can strengthenit directly without need of food or drink, 105 or quit it at will to vivify Divine Body and Sacred Circle of the Senses 155 that of another. 106 Experiencing the eternal satisfaction of his own uncreated nature, 107 the yogi is freed of disease and the infirmities of old age. 108 Inspired by this vitality, he has the inner strength to overcome the severest hardships and obstacles in the way of his ultimate goal. 109 The Stanzas on Vibration explain: Lassitude ravages the body and this arises from ignorance but if this is eliminated by an expansion of consciousness (unmesa), how can [ignorance], deprived of its cause, continue to exist? 110 Doubt (sanka) is the source of every spiritual ailment. 111 A

man in doubt about his true nature, and the path to follow by which he may come to realise it, is constantly overcome by difficulties. 112 "Due to ignorance" teaches the Sarvacaratantra, "the deluded man is in doubt and thus [suffers the cycle] of birth and death." 113 Caught between conflicting alternatives, the mind shifts from one to the other becoming more entangled in its own thoughts as it does so. Doubt thus contracts consciousness 114 sullying it with the turbulence of confused thinking. 115 At once both the root of "the ancient tree of transmigration" and the first sprout of its seed, 116 doubts deprives man of the innate bliss (sahajananda) of his own nature. 117 Like a thief it steals away the wealth of true knowledge and reduces man to an imagined state of poverty in which he feels hemmed in on all sides by constraints and limitations. 118 Ignorance and its sister, doubt, are the essence of all impurities (mala) which sully consciousness. To the degree in which the yogi is cured of this ailment, namely, the lassitude (glani) of doubt, his true nature becomes manifest, just as raw gold when heated is freed from dross. When the yogi recognises the all-powerful expansion (unmesa) of his consciousness, he is projected beyond the realm of relative distinctions between virtue and vice. Thus overcoming all doubt as to what he should or should not do, he penetrates into the pulsing Heart of Bhairava's consciousness. 119 As the Stanzas on Vibration say: An individual who [though] desirous of doing various things [but] is incapable of doing them due to his innate impurity, [experiences] the Supreme state when the disruption (ksobha) [of his ego] ceases. Then [the soul realises] that his [true] uncreated nature is [universal] agency and perceiving subjectivity and so he knows and does whatever [he] desires. 120 the individual soul is(=) a combination of (e) conscious (cit) and 156 The Doctrine of Vibration unconscious (acii) elements 121 brought together when he allows the power of his authentic nature to be obscured by (e) his own impurity. This impurity is (=) three-fold. The first is technically called (=) the 'Impurity of Individuality' (anavamald). Due to (e) this impurity, the individual soul fails to (e) recognise his all-embracing fullness (purnata) and is disturbed by (e) his craving

for (e) experience (bhoga) based on (e) the unreasoned assumption that he is (=) incomplete (apurnd) and in need of (e) something outside himself. 122 Thus it manifests as (=) desire or attachment (rdga). ]23 At the same time it is (=) primordial ignorance, 'the cause of the sprout of (e) transmigration'. 124 As such, it works in two ways. Firstly, it deprives (e) individual consciousness of freedom. The yogi affected by (e&eb) this impurity is incapable of (e) entering or emerging out of the states of contemplation as he wishes. Secondly, this impurity renders (eb) the individual unconscious of his freedom, which for(e) the yogi means that he is (=) liable to lose consciousness of (e) his true nature when he is (=) not in a state of (e) contemplative absorption. 125 The second impurity is (=) called the 'Impurity of Maya' (maylyamald). This comes into operation when consciousness has been contracted by (e) the Impurity of Individuality. It limits (e) the power of knowledge operating in (e&eb) the individual soul through (e&eb) the inner and outer organs of (e) perception. 126 Due to this impurity he perceives (e) diversity everywhere. 127 Finally, the third impurity, that of Karma, comes into play when (e) the individual, deprived of (e) the freedom and knowledge of (e) unconditioned consciousness, acts (e&eb) in his limited way prompted by (e) desires and fears for (e) his personal gain. Disturbed by (e) these three impurities the soul cannot find rest in (eb) himself.

And

$$E_6 > n_6$$

Replacing $E_6$ and $n_6$ in these conditions leads to the form

If $A_6^2 > 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6$.

If $A_6^2 < 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6 .+ (Y_6 - A_6^2/2)^2/2M_6(X_6 - Y_6)$.

Note that by normalizing $A_6$ by $M_6$ and $X_6$ and $Y_6$ by $W_6^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_6^2 + (M_6 + U_6)\lambda_6 - (M_6 + V_6)\lambda_6 e^{2T_6\lambda_6} + (M_6 U_6 + Y_6)$$

$$-[M_6(U_6 + V_6) - X_6]\, e^{2T_6\lambda_6} + M_6 V_6\, e^{2T_6\lambda_6} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_6 = \mathrm{i}\, Z_6$ *and* separate

(A3.1)into real and imaginary parts

$$Z_6^2 + (M_6 + V_6) Z_6 \sin(Z_6 T_6) - (U_6 M_6 + Y_6) + (M_6(U_6 + V_6) - X_6)$$

$$\mathrm{X}\cos(Z_6 T_6) - 2M_6 V_6 \cos^2(Z_6 T_6) + M_6 V_6 = 0$$

$$(M_6 + U_6) Z_6 - (M_6 + V_6) Z_6 \cos(Z_6 T_6) + (M_6(U_6 + V_6) - X)$$

$$\mathrm{X}\sin((Z_6 T_6) - M_6 V_6 \sin(Z_6 T_6) = 0$$

Using (A.3.3) to eliminate $Z_6$ and $Z_6^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_6 =$

$1 - \cos Z_6 T_6$ and $A_6 = U_6 - V_6$ we obtain

$$4M_6^2 V_6^2 W_6^4 + 2M_6 V_6 W_6^3[(M_6 - V_6)^2 + A_6(M_6 + V_6) + 2(A_6 M_6 - X_6)]$$

$$+ W_6^2 [M_6 - V_6)^2 (2 A_6 M_6 + Y_6 - X_6) + (2A_6 M_6 - X_6)^2 - A_6 (M_6 - V_6)(2A_6 M_6 - X_6)$$

$$+ M_6 V_6 (A_6^2 + 4(X_6 + Y_6))]$$

$$+ W_6 [A_6^2 + 2X_6)(2 A_6 M_6 - X_6) + 2A_6 Y_6 (M_6 + V_6) - M_6 A_6^3] + A_6^2 (X_6 M_6 - Y_6) = 0$$

Being a quartic in $W_6$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_6 > 0$$

$$and$$

$$C_6 > (X_6 - Y_6)/2 M_6$$

$recalling\ that\ A_6 = U_6$ and defining $C_6 = U_6 + V_6$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_6 (0, U_6 V_6 = A_6^2 (X_6 + Y_6)$$

$$Q_6 (2, U_6 V_6) = (C_6 + 2M_6^2)^2 (2C_6 M_6 + Y_6 - X_6)$$

So that given (A3.6) the extreme $W_6$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_6$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_6$ the existence of a D-position at this point is

determined by letting cos $(Z_6 T_6) = -1\ in\ (A3.3)$

ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 e^{-\tau_6\lambda_6} + X_6 + Y_6 = 0 \qquad (A1.1)$$

Where $V_6 = dB_6/dR_6$ $U_6 = dD_6/dR_6 + N_6 dG_6/dR_6$ $X_6 = N_6G_6dF_6/dR_6$ and $Y_6 = -N_6G_6dM_6/dR_6$ (evaluated at the equilibrium). Note that $U_6X_6$ and $Y_6$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_6$ as $(T_6)$ . (It follows that $f(T_6)$ 0 for all $T$, and $\int_0^\infty f(T_6)dt = 1.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 \int_0^\infty e^{-\tau_6\lambda_6} f(T_6) + X_6 + Y_6 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MODULE NUMBERED TWO**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed

the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).. Haimovici's and Nunny's models have been applied to some of the Deleuzean and Heideggerian problems too : Scoring on that count are :

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time

delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t - T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if

only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

TIME DELAY

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

## MATURATION AND CONSUMMATION OF CONFEDERATION OF DEBIT AND CREDIT TYPES:

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t-T_7)F_7(R_7(t-T_7) - N_7(t)M_7((R_7(t) \qquad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7} + \frac{N_7 dG_7}{dR_7} - \frac{dB_7}{dR_7} > \frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7} + \frac{dM_7}{dR_7}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes $(L_7)$ and the other derived from the non-delayed or "immediate" processes$(I_7)$:as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_7 = \frac{N_7 G_7}{2M_7}\frac{d(7nF_7)}{dR_7} (\geq 0) \qquad (7a)$$

$$I_7 = -\frac{N_7 G_7}{2M_7}\frac{d(7nM_7}{dR_7} \ (\geq 0) \qquad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7 > L_7 - I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_7 > L_7 - I_7, \quad \text{when} \quad P_7^2/4 > I \qquad (8a)$$

$$P_7 > L_7 - I_7 + \frac{(I_7 - P_7^2}{(L_7 - I)} \quad \text{when} \quad P_7^2/4 < I. \qquad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

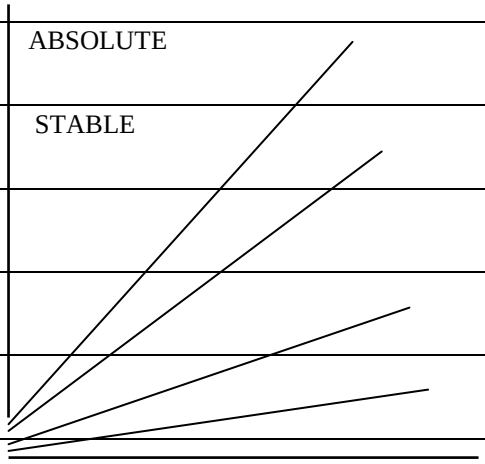



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit

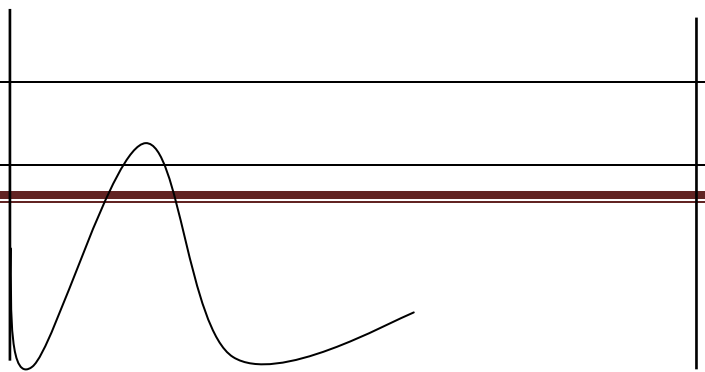
(like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

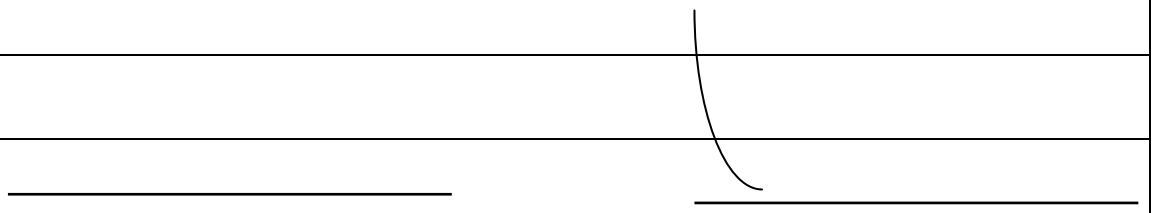
Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the





TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t-T_7)F_7\big(R_7(t-T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t-T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation $(dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7)$ and the effective prey regulation $(dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7)$ must be positive.

The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the

total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

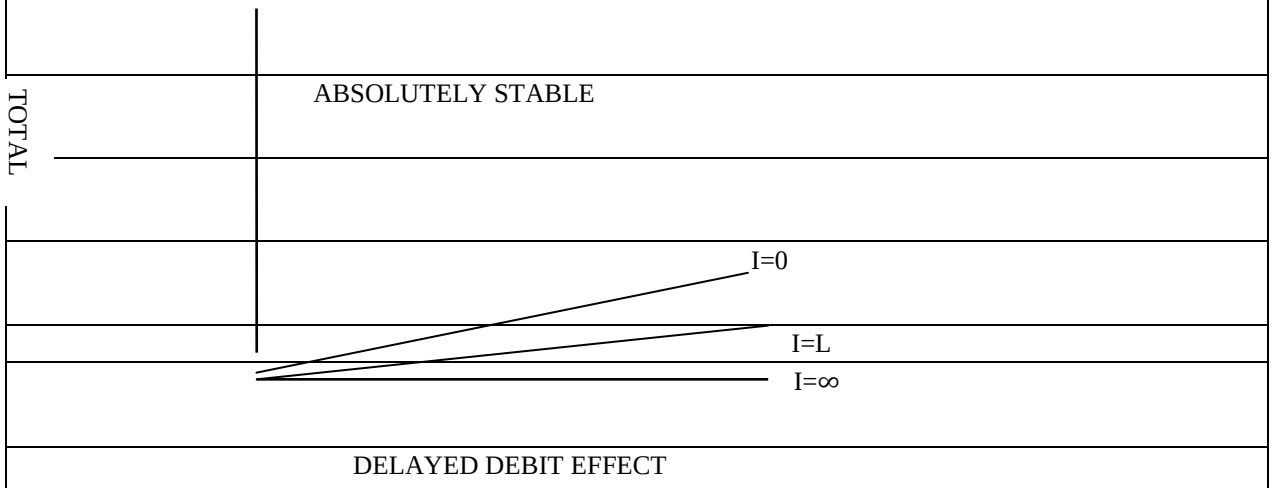



FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

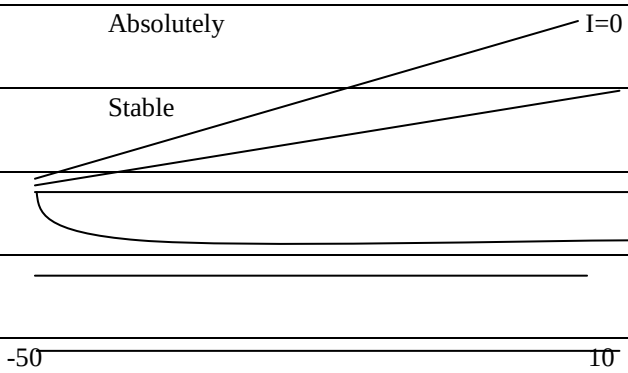
Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.

5

DEBIT REGULATION




**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay.

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^\infty \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7\left(U_7 - V_7\int_0^\infty cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7\right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7\int_0^\infty cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$then given the constraints of $f(T_7)$, $(A1.5)$*can never be satisfied Co*nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7\lambda_7 e^{-T_7\lambda_7} + [(U_7 - V_7)M_7 + Y_7] - [((U_7 -$$

V 7M 7−]$e^{-T7\lambda 7}$=0

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$Z_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7)\cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7)\sin(Z_7 T_7) = 0$$

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $Z_7^2$ and $Z_7$in (A2.2):

$W_7^2(X_7^2 - 3M_7 A_7 X_7 - M_7^2(X_7 - Y_7) + 2M_7^3 A_7 + 2A_7^2 M_7^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_7$=1-cos $(Z_7 T_7)$.

In the absence of a delay the system is stable for $A_7 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7 > (X_7 - Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$$A_7 > [\ X_7 - Y_7/2\ M_7$$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_7, A_7) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_7 \le 2$ *is* of interest First consider the extreme values

$$Q_7(0, A_7) = A_7^2\ (X_7 - Y_7)$$

$$Q_7(2, A_7) = (A_7 + 2\,M_7)^2\,(2A_7 M_7 - (X_7 - Y_7))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7 M_7^2 - 3M_7 X_7\right) + W_7\left(3M_7 A_7^2 - 2A_7 X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7+Y_7\right)$$

$=2M_7^2W_7^2(M_7+\mathrm{A}_7)+M_7A_7^2W_7+\left(2A_7M_7-X_7+Y_7\right)(A_7W_7+M_7W_7^2$

$+(2-W_7)((2M_7X_7+M_7Y_7)W_7+A_7(X_7+Y_7)),$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $W_7$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$has no minimum in the range $0 \le W_7 \le 2$.Defining

$Q_7\left(W_7,A_7\right)=a_7W_7^2+c_7 where\ a_7\ b_7\ and\ c_7$ Are functions of $A_7$we can define the condition for a minimum in the range 0-2

condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7 \ge b_7 \ge 0$$

And

$$b_7^2 \le 4a_7c_7$$

$$Case(i): Y_7=0$$

$Let\ Y_7=0$ And try the substitution A= $X_7=/2M_7\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}\ (W_7 X_7 / 2M_7)$ is positive for *W<2. Since* $Q_7$ increases with $A_7$(see A2.9), increasing $A_7$will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

$Case\ (ii): Y_7 = y_7\ X_7$ ,

*Let* $Y_7 = y_7\ X_7$And try the substitution $A_7 = X_7/(2 + A_7\ y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + n_7 y_7)$ we have

$$W_7^2\ y_7\left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2\left(2 - n_7 + n_7 y_7\right)\right]$$

$$- W_7\ X_7\left[\frac{X_7(1 + n_7 y_7)}{M_7^2(2 + n_7 y_7)^2} + 2\ y_7\left(n_7 - 1\right) + \frac{X_7^2(7 + y_7)}{M_7^2(2 + n_7 y_7)}\right. = 0.\ (A2.12)$$

Since $n_7$ and $A_7$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7), the\ condition (A2.77?) can\ be\ written\ as$

$[\frac{X_7(E_7 - 7)}{E_7^2}]^2[\frac{X_7^2}{M_7^2} - 4\ y_7\ E_7^2\frac{X_7}{M_7^2} + 4y_7(n_7 - 2)E_7^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$2\ E_7[y_7 E_7 + (y_7\ E_7(7 + y_7)(E_7 - n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - y_7 E_7(7 + y_7)(E_7 - n_7))^{7/2}]$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2$) is large or small then (A2.6) will have a real solution although if may not be in the range (0 W 7 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7 - 7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2 - 2y_7\frac{(E_7 + n_7^2 y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right) - 4y_7(E_7 - n_7)$$ 0

And has the solution

$\frac{X_7}{M_7^2} > \frac{E_7}{(E_7-7)}[y_7(E_7 + n_7^2 y_7) + [y_7^2 (E_7 + n_7^2 y_7)^2 + 4 y_7 (E_7 - 7)(E_7 - n_7)]^{7/2}]$

Given $E_7 > n_7$ this lower limit to $X_7 / M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7 / M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$ in the appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7 < n_7$ $(ie., 2A_7 M_7^2 < X_7 - Y_7)$ then (A2.7)and (A2.8)show that $Q_7$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - (y_7 E_7 (7 + y_7)(E_7 - n_7)]^{7/2}]$$

And

$$E_7 > n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2 > 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.

If $A_7^2 < 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.+ $(Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$ and $X_7$ and $Y_7$ by $W_7^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSILIDATAN**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7\lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]\, e^{2T_7\lambda_7} + M_7 V_7\, e^{2T_7\lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i}\, Z_7$ *and* separate

(A3.1)into real and imaginary parts

$$Z_7^2 + (M_7 + V_7) Z_7 \sin(Z_7 T_7) - (U_7 M_7 + Y_7) + (M_7(U_7 + V_7) - X_7)$$

$$\mathrm{X}\cos(Z_7 T_7) - 2M_7 V_7 \cos^2(Z_7 T_7) + M_7 V_7 = 0$$

$$(M_7 + U_7) Z_7 - (M_7 + V_7) Z_7 \cos(Z_7 T_7) + (M_7(U_7 + V_7) - X)$$

$$X \sin\left(\left(Z_7 T_7\right) - M_7 V_7 \sin\left(Z_7 T_7\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_7$ and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - cos\ Z_7\ T_7$ and $A_7 = U_7 - V_7$ we obtain

$$4 M_7^2 V_7^2 W_7^4 + 2 M_7 V_7\ W_7^3 [(M_7 - V_7)^2 + A_7 (M_7 + V_7) + 2(A_7 M_7 - X_7)]$$

$$+ W_7^2 [M_7 - V_7)^2 (2 A_7\ M_7 + Y_7 - X_7) + (2A_7\ M_7 - X_7)^2 - A_7 (M_7 - V_7)(2A_7\ M_7 - X_7)$$

$$+ M_7 V_7 (A_7^2 + 4(X_7 + Y_7))]$$

$$+ W_7 [A_7^2 + 2X_7)(2 A_7\ M_7 - X_7) + 2A_7 Y_7 (M_7 + V_7) - M_7 A_7^3] + A_7^2 (X_7\ M_7 - Y_7) = 0$$

Being a quartic in $W_7$ (A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7 > 0$$

$$and$$

$$C_7 > (X_7 - Y_7)/2\ M_7$$

$recalling\ that\ A_7 = U_7$ and defining $C_7 = U_7 + V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0, U_7 V_7 = A_7^2 (X_7 + Y_7)$$

$$Q_7\left(2, U_7 V_7\right) = \left(C_7 + 2M_7^2\right)^2 (2C_7 M_7 + Y_7 - X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1\ in\ (A3.3)$

## ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

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# Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price, University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss**

(389) In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them.

(390) Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors.

**Classical vs. relativistic black holes**

(391) Something like a black hole exists within (eb) Newton's classical theory of gravity.

(392) In that theory, an energy argument tells us that there is (=) an escape velocity vesc=2GM/R−−−−−−−√ from (e) the surface of any spherical object of (e) mass M and radius R

(393) If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity.

(394) Thus the condition for such an "unseeable" object is R<2GM/c2

(395) In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape.

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| (396) This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity. |
| (397) To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it. |
| (398) The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c . |
| (399) This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed. |
| (400) An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity. |
| (401) Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable. |

**Stationary black holes**

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| (402) In the Newtonian theory, gravity is described by (e) the potential $\Phi$. |
| (403) Inside a spherical object the form of (e) $\Phi(r)$ depends on (=) the interior structure, but in the vacuum outside (e) matter the potential $-GM/r$ depends only on (=) the interior mass. |
| (404) Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object. |
| (405) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward. |
| (406) This horizon radius $RH=2GM/c2$ is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory. |
| (407) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution. |
| (408) Radius RHhere actually means that the area of the event horizon is (=) $4\pi R2H$. |
| (409) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution. |
| (410) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity. |
| (411) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own. |
| (412) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole. |
| (413) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density. |
| (414) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed. |
| (415) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole. |

(416) In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter.

(417) Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum.

(418) In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior.

(419) The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object.

(420) Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum.

(421) These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.]

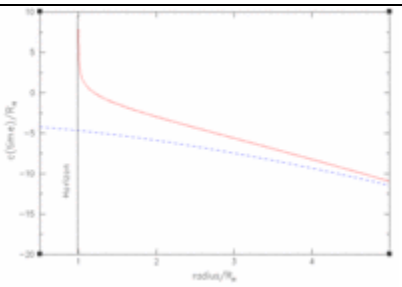
(422) All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time.

(423) But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?"

(424) The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point.

(425) One of these coordinates is (=) called the "coordinate time."

(426) Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate.



(427) Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole

(428) At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole.

(429) Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations.

(430) For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree

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| (431) An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events. |
| (432) The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving. |
| (433) Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle. |
| (434) The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return. |
| (435) But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon. |
| (436) Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future. |
| (437) The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?" |

**Black hole parameters**

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| (438) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin). |
| (439) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun. |
| (440) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits. |
| (441) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form. |
| (442) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds. |
| (443) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2. |
| (444) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed. |
| (445) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer. |
| (446) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2) |
| (447) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials. |
| (448) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable |
| (449) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite. |
| (450) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole. |

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| (451) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars. |
| (452) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity. |
| (453) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity. |
| (454) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite. |
| (455) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity. |
| (456) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum. |
| (457) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship. |
| (458) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= $(4\pi G2/c4)[(M+M2-a2\sqrt{\phantom{xx}})2+a2]$. |
| (459) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored. |

**Dynamical black holes**

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| (460) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape. |
| (461) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, **a horizon can be born**. |
| (462) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid. |
| (463) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward. |
| (464) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations. |
| (465) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes. |
| (466) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3. |
| (467) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/ |
| (468) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers. |
| (469) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission. |

(470) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.

(471) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.

(472) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.

(473) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.

(474) The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature $6\times10-8$ (Mo/M) K.

(475) Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.

**Astrophysical black holes**

(476) Black holes in our Universe can be grouped as: primordial black holes, stellar-mass black holes, and supermassive black holes. The first is highly speculative; the second and third are broadly accepted.

(477) Primordial black holes of all masses are postulated to have formed from (e) quantum fluctuations in the early Universe, but those with (e&eb) mass less than around 1013kg would have already evaporated due to Hawking radiation.

(478) No significant observational evidence has yet been found of (e) the existence of these objects [MacGibbon and Carr (1991)].

(479) Stellar-mass black holes, ranging from a few to a few tens of solar masses, are (=) normal but rare endpoints in (eb) the evolution of massive stars.

(480) When a star exhausts its nuclear fuel and cools, it must collapse unless (e) it is supported by nonthermal forces.

(481) It is known that such nonthermal forces cannot (e) resist gravitational compression for (e) masses greater than around 1.5 to 3Mo. (The uncertainty is due to (e) uncertainty in our knowledge of nuclear physics at high densities.)

(482) There are many stars far more massive than this, but in their death throes massive stars expel (eb) a great deal of their mass in supernova explosions

(483) Even stars initially as massive as 20-30Mo may blow off enough mass to (e) leave a remnant neutron star smaller than 1.5Mo.

(484) Stars more massive than 20-30Mo may form (eb) a neutron star core, only to have it collapse due to (e) fallback of material from (e) the stellar mantle. Still more massive stars (above ~40Mo) may form (eb) a black hole directly in collapse, with or without (e&eb) a supernova explosion.

(485) These masses are (=) rather uncertain at present, as our understanding of supernova explosion mechanisms is (=) still evolving [Woosley and Janka (2005), Muno (2006)].

(486) Collapsing gas clouds larger than ~100Mo are expected to dissipate from (e) radiation pressure, which would prevent (e) more massive stars from forming, and consequently set (eb) an approximate upper limit for stellar mass black holes.

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| (487) Observational evidence for stellar-mass black holes comes primarily from (e) X-ray astronomy. |
| (488) A black hole in a close binary orbit can form (eb) an accretion disk (see Figure 1) of matter pulled off (er) a normal stellar companion. In its inspiral, disk material is heated by shearing, and becomes (=) a strong X-ray emitter. |
| (489) If observations reveal a point-like X-ray emitter with (e&eb) a mass inferred from (e) orbital dynamics to be (=) above that possible for a neutron star, it becomes (=) a black hole candidate. |
| (490) t present there are over a dozen such black hole candidates, including the first object to be identified as (=) a black hole, the X-ray source Cygnus X-1 (estimated mass 10±3Mo) [Casares (2006)]. |
| (491) The other class of observationally-supported black holes is (=) that of supermassive black holes, ranging from hundreds to (e&eb) billions of solar masses. |
| (492) Evidence for the existence of black holes at (eb) the upper end of this range is overwhelming. |
| (493) By contrast, the existence of black holes with (e&*eb) mass roughly of order 102 to 104Mo, referred to as "intermediate-mass black holes," is speculative. |
| (494) The supermassive holes may have begun as (=) primordial or stellar-mass black holes, but have grown through (e&eb) absorption of stars or gas or through mergers with (e&eb) other holes [Ferrarese and Ford (2005)]. |
| Figure 4: Observed orbits of stars near (eb) the dark compact object at the center of (e) our Milky Way galaxy: still frame from an animation. Animation created by Prof. Andrea Ghez and her research team at UCLA, from data sets obtained with the W. M. Keck Telescopes: http://www.astro.ucla.edu/~ghezgroup/gc/ |
| Evidence for supermassive black holes originally came from (e) quasars, small but intense radio sources seen at (eb) cosmological distances. |
| Their huge luminosities implied (eb) high mass, while (e) their rapid variability implied (eb) extremely small size. |
| More recently, measurements of Doppler shifts of stars and gas in the centers of galaxies have shown (eb) that compact objects of mass greater than 106Mo reside in (eb) the cores of most galaxies; the very small size of these central objects rules out any plausible alternative to the black hole explanation. |
| The black hole at the center of our own Milky Way galaxy has (e) a measured mass of ~3×106Mo , based on (e) the orbits of stars near the radio source Sagittarius A* associated with (e&eb) the Galactic core [Ghez et al. (2005); see Fig. 4]. |
| (500) Black holes are of particular interest for (e) gravitational waves, and vice versa. |
| (501) For gravitationally-bound systems, the typical maximum gravitational wave amplitude (dimensionless strain) is (=) h~(2GM/c2)2/(RD)~R2H/(RD) , where D is the distance to the system, M and R are the mass and size of (e) the system, and RH is the horizon radius of (e) a black hole of the same mass. |
| (502) From this it is clear that strongest gravitational waves will involve (e&eb) systems at or near black hole compactness: in particular, systems containing (er) black holes in close proximity to one another [Flanagan and Hughes (2005)]. |
| (503) Conversely, a system containing only black holes would not be expected to radiate (eb) any sort of radiation other than gravitational waves, and in any system containing (e) black holes, gravitational waves provide (eb) the most accurate probes of conditions deep in (eb) their gravitational wells. |
| (504) In particular, gravitational wave observations could definitively test (e&eb) whether a compact mass is truly a black hole [Cutler and Thorne (2002)]. |

(505) The strongest intrinsic source of gravitational waves in the present Universe is (=) the collision of two comparable-massed black holes.

(506) A sufficiently tight black hole binary system will lose (e) energy through emission of (e&eb) gravitational waves, causing (eb) the orbit to circularize and then slowly shrink over time.

(507) Figure 5: Gravitational waves from (e) a black hole binary inspiral: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/

(508) When the orbital radius approaches (eb) the Schwarzschild radius of the system, complex nonlinear dynamics come into play (e&eb) and the final stages of (e) the inspiral and merger can only be modeled through (e&eb) numerical simulations.

(509) These simulations are quite challenging, but recent breakthroughs [Pretorius (2005)(See Fig. 5), Campanelli et al.(2006), Baker et al.(2006)] have led to (eb) the first complete inspiral-merger gravitational waveforms.

(510) After the black holes have coalesced, they form (eb) a single highly-distorted black hole which quickly settles down to (e) a quiescent Kerr state through (e&eb) mission of quasinormal "ringdown" gravitational waves [Pretorius (2005), Campanelli et al. (2006), Baker et al. (2006)].

(511) Another gravitational emission mechanism for supermassive black holes would come from (e) their capture of compact stellar objects, particularly stellar-mass black holes.

(512) Unlike comparable-mass mergers, these captures would be on (eb) highly eccentric orbits that would not circularize before merger, resulting in (eb) intricate rapidly precessing orbits in (eb) the final months or years of inspiral.

**(513)** The consequent gravitational-wave signal would allow (eb) for exquisitely precise measurements of the parameters of (e) the supermassive black hole, as well as tests of any deviation from (e) the predictions of general relativity [Cutler and Thorne (2002)].

**(514)** For stellar-mass black holes (tens of Mo), the final stages of inspiral and merger occur at (eb) frequencies of hundreds to thousands of Hz, a regime targeted by (e) ground-based gravitational-wave observatories.

**(515)** Systems of supermassive black holes that are (=) millions of Mo emit at frequencies of tens of (e) mHz, which will be detectable by (e) proposed space-based (e) gravitational-wave detectors.

**(516)** Black holes in the billions of Mo, or a stochastic background of signals from (e) more moderate supermassive black holes at (eb) earlier stages of their inspiral, might emit (eb) waves at nHz frequencies that would be measurable through (e&eb) years-long correlated timing measurements of radio pulsars [Jenet et al. (2005)].

**An equation to quantify the origins of life on other planets July 5, 2016 by Bob Yirka report**

(517) (Phys.org)—A pair of researchers, one with the Columbia Astrobiology Center in New York, the other with the University of Glasgow in the U.K. has come up with a mathematical equation that when (e) solved is meant to offer (eb) a means for estimating how often life begins on other planets. In their paper published in Proceedings of the National Academy of Sciences, Caleb Scharf and Leroy Cronin describe their equation, how they came up with it and why they believe it might become more useful as scientists learn more about the true nature of other planets and solar systems.

(518) Back in the 1960's, an astronomer named Frank Drake came up with a formula for (e) estimating the likely number of (e) alien civilizations that might be capable of transmitting radio signals in such a way as to be (=) recognizable by receivers here on earth.

(519) That equation, named quite naturally the **Drake Equation** has been the unofficial standard-bearer for half a century, despite it having no solution because (e) some of the parameters are still unknown.

(520) The same holds true for the new equation developed by (e) Scharf and Cronin, but the results given should the unknowns ever be discovered would offer a different sort of answer than Drake was looking for.

(521) To build their formula, the duo came up with several parameters to solve for (e) **abiogenesis**, which is the likelihood of an event occurring that leads to (eb) life beginning, they included (e) the number of possible building blocks, the mean number of such blocks per possible organism, the availability of building blocks that might exist during a given time period, expressed as a fraction and the probability that the existence of the building blocks would actually lead to (eb) life starting in a given unit of time. It looks like this: **Nabiogenesis (t)= Nb · 1/no · fc · Pa · t**

(522) In essence the formula is meant to suggest (eb) that the probability of life beginning on a given planet is very likely connected to (e&eb) whether there are building blocks available on a given planet, and how much of them there might be.

(523) As more is learned by space researchers it is hoped that the formula could help narrow down (e) search targets by offering (eb) a statistical probability of success for (e) a given planet.

(524) The researchers note that the building blocks available do not necessarily have to be (=) of the type that led to (eb) life beginning on our planet.

**(525)** They also suggest that implied in the formula is (=) the likelihood that life beginning events are (=) more likely to occur in solar systems where (e) there are multiple planets, which allow (eb) for an opportunity to share materials that could lead to (eb)building blocks **More information: Quantifying the origins of life on a planetary scale, Proceedings of the National Academy of Sciences, www.pnas.org/cgi/doi/10.1073/pnas.1523233113**

**Abstract**

(526) A simple, heuristic formula with (e&eb) parallels to the Drake Equation is introduced to help focus discussion on open questions for (e) the origins of life in a planetary context.

(527) This approach indicates (eb) a number of areas where (e) quantitative progress can be made on parameter estimation for (e) determining origins of life probabilities, based on (e) constraints from Bayesian approaches.

(528) Authors discuss a variety of "microscale" factors and their role in determining (eb) "macroscale" abiogenesis probabilities on (e&eb) suitable planets.

(529) Authors also propose that impact ejecta exchange between (e&eb) planets with parallel chemistries and (e&eb) chemical evolution could in principle amplify (eb) the development of (e) molecular complexity and (e&eb) abiogenesis probabilities.

(530) This amplification could be very significant, and both bias conclusions about abiogenesis probabilities based on (e) the Earth and provide (eb) a major source of variance in the probability of life arising in (eb) planetary systems.

**(531)** Authors use heuristic formula to suggest (eb) a number of observational routes for improving constraints on (e) origins of life probabilities. **Read more at: http://phys.org/news/2016-07-equation-quantify-life-planets.html#jCp**

**Brian Koberlein Shared publicly - Yesterday 8:00 PM the Illusion of Truth**

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| (532) | Physical theories are often presented as (=) a description of what's really going on. |
| (533) | Forces act on (e&eb) a baseball, causing (eb) baseball to fall. |
| (534) | Atoms collide and fuse in (eb) the heart of a star, releasing heat and energy. Science is true, as is often said, and our scientific theories encapsulate this truth. But this isn't entirely true. |
| (535) | For example, quantum mechanics is a strange, sometimes confusing theory. Objects can be (=) particles and waves. |
| (536) | They do strange things like tunnel through (e&eb) barriers, and can even appear and disappear (e&eb). |
| (537) | Given all this strangeness, what's really going on? What's the true nature of quantum reality? It depends on (=) which approach you want to take. |
| (538) | In one view quantum objects are described by (e) a wave-like probability function. |
| (539) | When you interact with (e&eb) or "observe" (e) the object this "wavefunction" collapses into (eb) a definite state. |
| (540) | That's view commonly presented, but a quantum system can also be described by (e) its transitions between energy states. |
| (541) | Since this uses a mathematical method involving (e&eb) matrices, it's known as (=) matrix mechanics. |
| (542) | Both wavefunctions and matrix mechanics give (eb) the same results, but their view of what really goes on is very different. |
| (543) | Then there's the path integral method |
| (544) | . Rather than a wavefunction or matrix transitions, path integrals imagine quantum objects can take almost any path between (e&eb) two states. |
| (545) | By summing all the possible paths you can derive (eb) the odds that it will occur. |
| (546) | So which one is true? Are quantum objects distributed waves of (e) probability? |
| (547) | Are they simply transitions between (e&eb) energy states? |
| (548) | Do they take an infinite number of (e) paths between states? |
| (549) | Each of these models makes (eb) the same predictions, so one version is no more "true" than the others. |
| (550) | What happens in practice is that we'll use whatever method is useful at the time. They are equivalent models, so the best model for the job is the one we'll use. The only reason the wavefunction view is so common is that it's (=) the version usually taught to introductory students. You might think this uncertainty of truth is due to (e) the behavior of quantum physics itself. It's so strange and counterintuitive that we can't (e) wrap our puny brains around what's really going on. But the same thing occurs in lots of other fields. Even something as straight forward as basic Newtonian physics |
| (551) | Toss a baseball in the air and (eb) Earth's gravitational force pulls it down. |
| (552) | The force of gravity is a simple truth, right? While we often describe classical motion in terms of (e&eb) forces and acceleration, we can also describe it in terms of (e&eb) energy and momentum. |
| (553) | In the Lagrangian and Hamiltonian approach, the path of a baseball is (=) the optimized path among possibilities. |
| (554) | In this view a baseball's path is (=) the extrema of an energy equation, and force can be derived as a necessary consequence of (e) this. |
| (555) | In the relativistic view the baseball follows (e&eb) a geodesic, which is (=) the minimal path through (e&eb) space and time. |

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| (556) So is a baseball's motion due to (e) a gravitational force, an energy extrema, or a spacetime geodesic? |
| (557) As with quantum theory, different approaches yield (eb) the same result. They are mathematically equivalent, so we can use whatever method is most useful at the time. |
| (558) At its core, science is less about truth and more about models. The metaphysics underlying a model is useful only as far as it allows us (eb) to make better predictions, generate (eb) new ideas, or bring models together as (=) a cohesive whole |
| (559) This is why we have no problem using (e) classical gravity to calculate the path of (e) a spacecraft through (e&eb) the solar system, while (e) using special relativity to account for (e) the Doppler shift of the spacecraft's radio signals. |
| (560) It's why we can use (e) quantum physics to study atoms in the morning, and general relativity to (e) study black holes in the afternoon. |
| (561) In regimes where models conflict with each other it isn't a failure of truth, but instead shows an opportunity to develop a better model. |
| (562) It could be that with each better model we move closer to (e&eb) the truth about reality. |
| (563) The truth is out there, and science strives to move towards (e&eb) that truth. |
| (564) It's a common view, and certainly the search for truth has driven (eb) many scientists to develop better and better models. |
| (565) But the real power of science is the recognition that what we have are models. Our models can be powerful, but they are always a bit tentative. There's always a chance that they might just be an illusion of truth. |

**Synopsis: A Single-Photon Cheshire Cat July 1, 2016 Researchers detected the polarization of a photon separate from the photon itself, just as the grin of Lewis Carroll's Cheshire cat can appear apart from the cat's body Synopsis figure APS/Carin Cain**

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| (566) In a quantum Cheshire cat, one property of a particle is present detached from (e) the particle itself, just as (=) the grin of the feline in Alice in Wonderland can exist in the absence of (e) the whole animal. |
| (567) Now a research team has successfully separated (e&eb) a single photon from (e) its polarization. |
| (568) The effect was previously demonstrated (eb) with neutrons and also with light in a classical scenario, but the latest experiment uses single photons. |
| (569) To produce the Cheshire cat, **Maximilian Schlosshauer and his colleagues** at the University of Portland in Oregon used an interferometer where (e) light hits a beam splitter, goes in two, opposite directions around a circuit (but along slightly offset routes), and then recombines at (eb) the entrance. |
| (570) The incoming polarization of light was horizontal, but this was switched to (e&eb) vertical for the counterclockwise loop. |
| (571) Each path included a device to adjust (e&eb) the polarization further and another to (e&eb) scatter away a small number of photons of (e) a designated polarization, the equivalent of detecting their presence. |
| (572) The team fired single photons into (e&eb) the interferometer so that only one photon was present at (eb) a time, traversing both paths simultaneously. |
| (573) They placed a horizontal polarizer in front of the output photon detector and gradually varied (e&eb) the difference between the two path lengths. |

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| (574) This variation generated (eb) a sine-wave interference pattern from (e) the exiting photons only when the polarization was tweaked in (eb) the counterclockwise path. |
| (575) At the same time, the average number of photons counted decreased only when (e) light was scattered from the clockwise path. |
| (576) Simultaneously demonstrating these different effects for two routes of a single photon constitutes (eb) evidence of the Cheshire cat. |
| **(577)** The effect might someday improve precision experiments if (e) an unwanted property could be moved away from the property being measured. **This research is published in Physical Review A. – David Ehrenstein David Ehrenstein is the Focus Editor for Physics.** |

**Observation of the quantum paradox of separation of a single photon from one of its properties James M. Ashby, Peter D. Schwarz, and Maximilian Schlosshauer Phys. Rev. A 94, 012102 – Published 1 July 2016**

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| (578) Authors report an experimental realization of the quantum paradox of (e) the separation of a single photon from (e) one of its properties (the so-called "quantum Cheshire cat"). |
| (579) Authors use a modified Sagnac interferometer with (e&eb) displaced paths to produce (eb) appropriately pre- and postselected states of (e) heralded single photons. |
| (580) Weak measurements of photon presence and (e&eb) circular polarization are performed in each arm of (e) the interferometer by introducing (eb) weak absorbers and small polarization rotations and analyzing changes in (eb) the postselected signal. |
| (581) The absorber is found to have (e) an appreciable effect only in one arm of the interferometer, while (e) the polarization rotation significantly affects (e&eb) the signal only when (e) performed in the other arm. |
| (582) Authors carry out both sequential and simultaneous weak measurements and find (eb) good agreement between measured and predicted weak values. |
| (583) In the language of Aharonov et al. and in the sense of the ensemble averages described by (e) weak values, the experiment establishes (eb) the separation of a particle from one its properties during the passage through (e&eb) the interferometer. |

**Viewpoint: Relativity Gets Thorough Vetting from LIGO Frans Pretorius, Department of Physics, Princeton University, Princeton, NJ 08544, USA May 31, 2016• Physics 9, 52 The gravitational-wave signal observed by the LIGO detectors shows no deviation from what general relativity predicts. Figure caption expand figure Frans Pretorius, APS/Carin Cain**

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| (584) Figure 1: The signal from one of the LIGO detectors in Hanford, Washington, is shown with two representations of (e) the best-fit numerical relativity (NR) waveform. |
| (585) The filtered NR waveform illustrates how the raw waveform is perceived by (e) the detector, |
| (586) On September 14, 2015, the LIGO interferometers made the first direct measurement of (e) gravitational waves [1]. |
| (587) The historic discovery came nearly 100 years after Einstein showed that his general relativity theory predicts (eb) the existence of gravitational waves and over 50 years since (e) the first quest to observe the Universe in gravitational waves [2]. |
| (588) The observed event, labeled GW150914, was (=) the collision of two black holes, one weighing 29 solar masses, the other 36 solar masses. |

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| (589) This binary coalesced into (e&eb) a remnant black hole weighing 62 solar masses and rotating (e&eb) at 67% the maximum rate allowed (eb) for a black hole. |
| (590) The most statistically significant portion of (e) the signal came from (e) the last several gravitational-wave cycles of the collision, where (e) the most nonlinear, strong-field dynamics of the theory manifest. |
| (591) One couldn't have asked for a better first event to learn about (e&eb) this regime of general relativity. |
| (592) The first thing that GW150914 can teach us is how closely nature follows (eb) Einstein's prescriptions for (e&eb) colliding black holes. |
| (593) Researchers from the LIGO team and the Virgo team have performed a series of tests to see if (e) any violations of (e) general relativity might be lurking in the data of GW150914 [3]. |
| (594) They first determined (eb) the theoretical waveform, or "template," from (e) general relativity simulations that best fit the measured signal (see Fig. 1). |
| (595) Their main finding is that the residual signal—obtained by (e) subtracting this template from (e) the data—is consistent with (e&eb) noise. |
| (596) This is, in a sense, the "everything" test, as (e) the template folds in (e&eb) all that general relativity predicts about the merger: the decay of the last few orbits of (e) the two black holes, their violent collision (e&eb) and astonishingly rapid relaxation to (e) a quiescent rotating black hole, the propagation of (e) gravitational waves through (e&eb) the cosmos, and finally the transformation of (e) these waves into(e&eb) a chirp in the LIGO detectors. |
| (597) Given the large signal-to-noise ratio of 24, the close agreement between the template and (=) the data puts limits on (e) alternative models. |
| (598) Specifically, the predictions of these models can only disagree with (e&eb) the general relativity prediction for GW150914 by, at most, 4%. |
| (599) In addition to this comprehensive test, the LIGO/Virgo teams also performed other checks that pick apart the signal, in order to address particular aspects of general relativity and black hole physics. In one such test, they compared the early and (e&eb) late parts of the waveform. |
| (600) From the waves generated during (e&eb) the initial inspiral, one can infer (eb) the masses and spins of the binary. |
| (601) Given these numbers, general relativity allows (eb) the mass and spin of (e) the remnant black hole to be calculated. |
| (602) The researchers compared this result to (e&eb) an independent measurement of (e) the mass and spin obtained from (e) the late-time ringdown, and to within (eb) uncertainties, the two estimates were (=) consistent. |
| (603) Another test isolated coefficients in the so-called Post Newtonian (PN) series describing (eb) the inspiral. |
| (604) The PN series is (=) an expansion of the gravitational-wave signal in powers of (e) the velocity vv of the sources relative to (e&eb) the speed of light cc. |
| (605) Observations of binary pulsars, which show (eb) a decrease in their orbital period due to (e) gravitational-wave emission, have constrained (e) PN terms, especially at low order. |
| (606) For higher-order terms associated with (e&eb) strong-field emission, the LIGO/Virgo teams were able to improve (eb) the limits by many orders of magnitude compared to (e&eb) the binary pulsar constraints. |

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| (607) At first glance it might seem these new results simply add to (e) the litany of previous tests that general relativity passed with flying colors. |
| (608) However, GW150914 is vastly more significant than that. Other tests have shown that general relativity is valid in (eb)the weak-field regime, but GW150914 gives (eb) the first direct observational evidence that general relativity is the correct theory of nature governing (e&eb) the extreme spacetime dynamics around the event horizon of (e) a black hole. |
| (609) One could argue that binary pulsar observations already told us that gravitational waves exist [4], but GW150914 is more luminous in gravitational waves than pulsar binaries by (e) at least 24 orders of magnitude. |
| (610) It's true that the evolution of the Universe as a whole is driven by (e) strong gravity, but it would be difficult to argue that precision cosmological measurements can constitute (eb) a test of general relativity when (e) one of the dominant inputs, dark energy, is still a mystery. |
| (611) As far as black holes are concerned, astronomers have (e) observed ultra compact dark objects that are consistent with (e&eb) black holes. |
| (612) The most striking data comes from (e) stellar orbits about the 4 million solar mass objects in (eb) the center of our Galaxy [5], but present observations cannot constrain it to be (=) a black hole as described by general relativity. |
| (613) GW150914 has given us the first morsels of quantifiable evidence that the rich physics of black holes is as predicted by general relativity. For example, the LIGO observations are consistent with (e&eb) so-called cosmic censorship [6], which says that gravitational singularities must remain hidden behind an (eb) event horizon. |
| (614) If a "naked" singularity were produced in (eb) the energetic collision of GW150914, the data would likely have shown (eb) some sort of anomaly, but in exactly what form, we don't know. |
| (615) Another example involves the "no-hair" theorem [7], which says that the only properties of a black hole are (=) its mass, spin, and charge (though astrophysical black holes are expected to have negligible charge). |
| (616) By consequence, these properties should determine (eb) the frequency and decay constants of (e) all the infinitely many ringdown modes of (e) a black hole. |
| (617) With GW150914, only the least-damped ringdown mode is (=) visible above the noise, but when combined with (e&eb) the remnant mass and spin inferred from (e) the inspiral, this is enough to count as (=) the first piece of evidence for a black hole's baldness. |
| (618) GW150914 should also be a test for (e) alternative theories of gravity or exotic alternatives to (e) black holes within general relativity, such as boson stars [8]. |
| (619) Some of these alternative models are concerned with (e&eb) explaining dark energy or dark matter and do not (e) therefore, predict any deviation from relativity at (eb) the scale of GW150914. |
| (620) However, other models do in fact deviate from (e) conventional relativity in their predictions of (e) the gravitational-wave signal from (e) a binary merger. |
| (621) The largest deviations should presumably occur during (e&eb) the last several gravitational-wave cycles, where (e) the signal is strong, and deviations should be (=) easy to detect or constrain. |
| (622) Embarrassingly though, the majority of alternatives that claim deviations can't (e) really say much about this event, as the two-body problem hasn't been solved in these theories. |
| (623) These models can only reproduce (eb) a fraction of the full gravitational-wave signal, (eb) so the strong constraints one might have hoped for (e) cannot be leveled. |

(624) Exotic alternatives within (eb) general relativity are in a similar boat, or worse for (e) objects like gravastars, which have (e) no well-defined theoretical framework to say what happens when (e) two such objects collide.

(625) Regardless, what is clear from the data is that any non-black-hole solution must imagine a form of (e) exotic matter with a viscosity larger than (e) any known material to explain (eb) the rapid damping of oscillations observed during the ringdown.

(626) The inferred event rate for heavy black hole mergers like GW150194 is high, and LIGO has another factor of 3 improvement in sensitivity slated before it reaches (eb) design sensitivity.

(627) Within 5 years then, we can expect dozens of detected events, with the loudest having a signal-to-noise ratio upwards of (e) 100.

(628) This will allow (eb) for exquisite insight into the nature of extreme gravity beyond (e) what GW150194 has already gifted us with. **This research is published in Physical Review Letters**.

**Tests of General Relativity with GW150914 B. P. Abbott et al. (LIGO Scientific and Virgo Collaborations) Phys. Rev. Lett 116, 221101 – Published 31 May 2016**

(629) The LIGO detection of GW150914 provides (eb) an unprecedented opportunity to study the two-body motion of (e) a compact-object binary in (eb) the large-velocity, highly nonlinear regime, and to witness the final merger of (e) the binary and the excitation of uniquely relativistic modes of (e) the gravitational field.

(630) Authors carry out several investigations to determine whether GW150914 is consistent with (e&eb) a binary black-hole merger in general relativity.

(631) They find that the final remnant's mass and spin, as determined from (e) the low-frequency (inspiral) and (e&eb) high-frequency (postinspiral) phases of the signal, are mutually consistent with (e&eb) the binary black-hole solution in general relativity.

(632) Furthermore, the data following the peak of GW150914 are consistent with (e&eb) the least-damped quasinormal mode inferred from (e) the mass and spin of the remnant black hole.

(633) By using waveform models that allow (eb) for parametrized general-relativity violations during (e&eb) the inspiral and merger phases, authors perform quantitative tests on (e&eb) the gravitational-wave phase in the dynamical regime and they determine the first empirical bounds on (e&eb) several high-order post-Newtonian coefficients.

(634) Authors constrain (e) the graviton Compton wavelength, assuming (eb) that gravitons are dispersed in (eb) vacuum in the same way as particles with mass, obtaining a 90%-confidence lower bound of 1013 km.

(635) In conclusion, within our statistical uncertainties, we find no evidence for violations of (e) general relativity in the genuinely strong-field regime of gravity. **Received 26 March 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.221101 © 2016 American Physical Society**

**Focus: Background Noise of Gravitational Waves March 31, 2016• Physics 9, 33**

(636) Assuming LIGO's recently detected black hole merger is not unusual; researchers revised upward their estimate of (e) the strength of the background noise coming from (e) distant mergers across the Universe. Figure caption expand figure Caltech/MIT/LIGO Lab (deleted due to spatial constraints)

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| (637) | Catching some waves The LIGO-Livingston interferometer consists of (e) two 4-km-long arms where (e) laser light bounces back and forth. Interference between the lasers from the two arms can signal tiny changes in (e&eb) the lengths of the arms resulting from gravi... |
| (638) | The recent detection of gravitational waves emitted from (e) a pair of merging black holes implies (eb) the existence of a large population of similar pairs, which should be generating (eb) a background of waves too faint to be detected with (e&eb) present technology. |
| (639) | Using this information, researchers now estimate that this background could be (=) roughly 10 times stronger than previously expected and might be found by (e) sensitive instruments in only a few years. |
| (640) | The background signal could help astrophysicists better understand how black hole binaries form and also pave the way to detecting the waves emitted by (e) the expanding Universe shortly after the big bang. |
| (641) | Gravitational waves were detected for the first time on 14 September 2015 by the Laser Interferometer Gravitational-Wave Observatory (LIGO). |
| (642) | By reflecting laser light between mirrors set 4 km apart, the team was able to detect (e&eb) tiny fluctuations in the distance between the mirrors as (=) the waves passed. |
| (643) | The waves were triggered by (e) an event, called GW150914; involving (e&eb) the merger of a pair of black holes having (e) masses of about 29 and 36 times that of the Sun. |
| (644) | Researchers expect that similar pairs of black holes, as well as neutron stars, should be merging constantly, creating (eb) gravitational waves that should bombard (e&eb) the Earth from all directions and make the equivalent of (=) a faint popping noise in a detector. |
| (645) | Previous estimates of the likely strength of this background had to make assumptions about (e&eb) how many pairs there might be and about the distribution of their masses. |
| (646) | These estimates remained speculative, as no actual binaries were known. |
| (647) | "What the recent observation has changed," says Nelson Christensen of Carleton College in Minnesota and the LIGO Scientific Collaboration (LSC), "is that we now know there really is (=) a population of black hole binaries. We have a small bit of information about it." |
| (648) | The new analysis—carried out by LSC researchers working with (e&eb) the European Virgo collaboration—estimates (eb) the energy density in the gravitational wave background in (eb) the frequency range below 100 Hz, where (e) the LIGO detector is most sensitive. |
| (649) | To make this estimate, they assumed (eb) a population of black hole binaries of various sizes, having (e) a range of distances from Earth, and created through (e&eb) a variety of scenarios. |
| (650) | They then calculated the amount of gravitational radiation that would result from (e) the merging of these objects. |
| (651) | The team then adjusted this estimate by using (e) details of the GW150914 event—especially the masses of the black holes involved—to make (eb) a most likely guess at the actual distribution of black hole binaries. |
| (652) | The idea is that the observed merger is (=) most likely representative of (e) the general population, rather than being extremely unusual. |
| (653) | If the black hole binary population is represented as (=) a bell curve, or Gaussian, the event should lie within (eb) the main bulge, not in the tails. |
| (654) | The observed mass of GW150914 suggests (eb) that the main population of black hole binaries could be (=) roughly 20 times more massive than expected, which would generate (eb) a background roughly 10 times "louder" than (e) previous estimates. |

(655) However, Christensen emphasizes that the result has (e) lots of uncertainty, as it's based on (e) only one event.

(656) "This estimate seems very plausible to me," says astrophysicist Maura McLaughlin of the University of West Virginia in Morgantown. She is a member of the NANOGRAV collaboration, which tries to detect gravitational waves at much lower frequencies than LIGO can observe. "I can't see any holes in their arguments. Unless LIGO got really lucky, then there must be many more of these systems than previously thought." This level of gravitational background implies (eb) that the LIGO and Virgo detectors might be sensitive enough to see it in 2–5 years, sooner than physicists previously expected, says Christensen, which has led to much anticipation. "Everyone is super excited," he says.

(657) Detecting the background would provide (eb) a confirmation of the presence of gravitational waves carrying (e&eb) energy throughout the Universe. It would also be useful, the researchers suggest, in gaining (e) further understanding of precisely how black holes form.

(658) And subtracting this background might someday allow (eb) detection of the fainter waves coming from (e) the expansion of the Universe moments after the big bang. This research is published in Physical Review Letters. –Mark Buchanan Mark Buchanan is a freelance science writer in Normandy, France.

**GW150914: Implications for the Stochastic Gravitational-Wave Background from Binary Black Holes B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 131102 – Published 31 March 2016**

(659) The LIGO detection of the gravitational wave transient GW150914, from (e) the inspiral and merger of (e) two black holes with masses $\gtrsim 30 M_{\odot}$, suggests (eb) a population of binary black holes with (e&eb) relatively high mass.

(660) This observation implies (eb) that the stochastic gravitational-wave background from binary black holes, created from (e) the incoherent superposition of (e) all the merging binaries in the Universe, could be (=) higher than previously expected.

(661) Using the properties of GW150914, authors estimate the energy density of (e) such a background from (e) binary black holes.

(662) In the most sensitive part of the Advanced LIGO and Advanced Virgo band for (e) stochastic backgrounds (near 25 Hz), authors predict $\Omega_{GW}(f=25\ \mathrm{Hz}) = 1.1^{+2.7}_{-0.9} \times 10^{-9}$ with 90% confidence.

(663) This prediction is robustly demonstrated for (e) a variety of formation scenarios with (e&eb) different parameters.

(664) The differences between models are (=) small compared to the statistical uncertainty arising from (e) the currently poorly constrained (e) local coalescence rate.

(665) Authors conclude that this background is potentially measurable by (e) the Advanced LIGO and Advanced Virgo detectors operating at their projected final sensitivity. DOI:http://dx.doi.org/10.1103/PhysRevLett.116.131102 © 2016 American Physical Society

**Study opens new realms of light-matter interaction July 15, 2016 by David L. Chandler. Some sentences might have been omitted due to spatial constraints. Please bear with me.**

(666) Emission spectra are a widely used method for (e) identifying chemical compounds; the bright lines reveal (eb) the different frequencies of light that can be emitted by (e) an atom.

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| (667) Here, a normal emission spectrum for (e) an atom in a high-energy state (top) is compared to (e&eb) the emission from the same atom placed just a few nanometers (billionths of a meter) away from (e) graphene that has been doped with (e&eb) charge carriers (bottom). |
| (668) For each energy-level transition, an orange line (or purple cloud) appears if (e) that transition is estimated to be faster than (e) one per microsecond — making it frequent enough to be (=) observed. **Credit: Massachusetts Institute of Technology** |
| (669) A new MIT study could open up new areas of technology based on (e) types of light emission that had been thought to be (=) "forbidden," or at least so unlikely as to be (=) practically unattainable. |
| (670) The new approach, the researchers say, could cause (eb) certain kinds of interactions between light and (e&eb) matter, which would normally take billions of years to happen, to take place instead (e&eb) within billionths of a second, under (e&eb) certain special conditions. The findings, based on a theoretical analysis, are reported today in the journal Science in a paper by MIT doctoral student Nicholas Rivera, Department of Physics Professor Marin Soljačić, Francis Wright Davis Professor of Physics John Joannopoulos, and postdocs Ido Kaminer and Bo Zhen. |
| (671) Interactions between light and matter, described by (e) the laws of quantum electrodynamics, are the basis of a (e) wide range of technologies, including (e) lasers, LEDs, and atomic clocks. |
| (672) But from a theoretical standpoint, "Most light-matter interaction processes are 'forbidden' by (e) electronic selection rules, which limits (e) the number of transitions between energy levels we have access to," Soljačić explains. |
| (673) For example, spectrograms, which are used to (e) analyze the elemental composition of materials, show (eb) a few bright lines against (e&eb) a mostly dark background. |
| (674) The bright lines represent (eb) the specific "allowed" energy level transitions in (eb) the atoms of that element that can be accompanied by (e) the release of a photon (a particle of light). |
| (675) In the dark regions, which make up most of the spectrum, emission at those energy levels is (=) "forbidden." |
| (676) With this new study, Kaminer says, "we demonstrate theoretically that these constraints can be lifted" using (e) confined waves within (e&eb) atomically thin, 2-D materials. |
| (677) "We show that some of the transitions which normally take the age of the universe to happen could be made to happen within (eb) nanoseconds. |
| (678) Because of this, many of the dark regions of a spectrogram become (=) bright once an atom is (=) placed near a 2-D material." |
| (679) Electrons in an atom have (e) discrete energy levels, and when they hop from one level to another they give off (eb) a photon of light, a process called spontaneous emission. |
| (680) But the atom itself is much smaller than the wavelength of the light that gets emitted—about 1/1,000 to 1/10,000 as big—substantially impairing (e&eb) the interactions between the two. |
| (681) The trick is, in effect, to "shrink" (e) the light so it better matches the scale of the atom, as (=) the researchers show in their study. |
| (682) The key to enabling a whole range of interactions, specifically transitions in (eb) atomic states that relate to absorbing or emitting light, is (=) the use of a two-dimensional material called graphene, in which light can interact with (e&eb) matter in the form of plasmons, a type of electromagnetic oscillation in the material. |
| (683) These plasmons, which resemble photons but have wavelengths hundreds of times shorter, are (=) very narrowly confined in the graphene, in a way that makes some kinds of interactions with (e&eb) that matter many orders of magnitude more likely than (e) they would be in ordinary materials. |

(684) This enables a variety of phenomena normally considered unattainable, such as (e) the simultaneous emission of multiple plasmons, or two-step light-emitting transitions between (e&eb) energy levels, the team says.

(685) This method can enable the simultaneous emission of (e) two photons that are "entangled," meaning (eb) they share the same quantum state even when (e) separated.

(686) Such generation of entangled photons is (=) an important element in quantum devices, such as those that might be used for cryptography.

(687) Making use of these forbidden transitions could open up (eb) the ability to tailor the optical properties of materials in ways that had not been thought possible, Rivera says.

(688) "By altering these rules" about the relationship between light and (e&eb) matter, "it can open (eb) new doors to reshaping (e&eb) the optical properties of materials."

(689) Kaminer predicts that this work "will serve as (=) a founding piece for (e) the next generation of studies on light-matter interactions" and could lead to (eb) "further theoretical and experimental advances in many fields which rely on (e) light-matter interactions, including (e) atomic, molecular and optical physics, photonics, chemistry, optoelectronics, and many others."

(690) Beyond its scientific implications, he says, "this study has possible applications across multiple disciplines, since (e) in principle it has potential to enable the full use of (e) the periodic table for optical applications."

(691) This could potentially lead to (eb) applications in spectroscopy and sensing devices, ultrathin solar cells, and new kinds of materials to absorb solar energy, organic LEDs with higher efficiencies, and photon sources for (e) possible quantum computing devices."From the standpoint of fundamental science, this work lays the groundwork for a subfield that just a few years ago was difficult to imagine and until now was largely unexplored," Soljačić says.

(692) "Two-dimensional materials confine (e) fields to a surface and motion to a plane, making possible (eb) many effects that are orders-of-magnitude too weak to (e) appear in a bulk volume," says Jason Fleischer, an associate professor of electrical engineering at Princeton University, who was not involved in this research.

(693) This work, he says, "systematically explores how 2-D materials improve light-matter interactions, laying a theoretical foundation for faster electronic transitions, enhanced sensing, and better emission, including the compact generation of broadband and quantum light." More information: N. Rivera et al: Shrinking light to allow (eb) forbidden transitions on the atomic scale, Science (2016). DOI: 10.1126/science.aaf6308 Journal reference: Science search and more info website: Provided by: Massachusetts Institute of Technology search and more info website. This story is republished courtesy of MIT News (web.mit.edu/news office/), a popular site that covers news about MIT research, innovation and teaching.

**Viewpoint: Peering Back to the Beginning of Time Lawrence M. Krauss, Department of Physics and School of Earth and Space Exploration, Arizona State University, Tempe, AZ 85287, USA and Mount Stromlo Observatory, Research School of Astronomy and Astrophysics, Australian National University, Weston, ACT 2611, Australia June 19, 2014• Physics 7, 64**

(694) The BICEP2 collaboration reports the detection of BB-mode polarization of (e) the cosmic microwave background—a signal that might originate from (e) gravitational waves created by (e) inflation during the very earliest moments in the evolution of our Universe. Figure caption expand figure P. A. R. Ade et al. [1]; Image on homepage: BICEP2 Collaboration (deleted due to restrictions on space)

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| (695) | Figure 1: (Top) BB-mode polarization pattern observed by (e) the BICEP2 telescope [1]. |
| (696) | The twisting pattern in the polarization of (e) the cosmic microwave background (CMB) might originate from (e) gravitational waves created by inflation in the very first moments of the universe. |
| (697) | Inflation—the hypothesis that the Universe underwent (e&eb) a phase of superluminal expansion in a brief period following the big bang—has (e) the potential of explaining, from first principles, why the Universe has (e) the structure we see today. |
| (698) | It could also solve outstanding puzzles of standard big-bang cosmology, such as why the Universe is, to a very good approximation, (=) flat and isotropic (i.e., it looks the same in all directions). |
| (699) | Yet we do not yet have a compelling model, based on (e) fundamental particle physics principles, that explains inflation. |
| (700) | And despite its explanatory power and a great deal of suggestive evidence, we still lack an unambiguous and direct probe of (e) inflation. |
| (701) | Theorists have developed different models for inflation, which all share (e&eb) a common, robust prediction: Inflation would have created (eb) a background of gravitational waves that could have (eb) an observable effect. |
| (702) | These waves would cause (eb) subtle, characteristic distortions of (e) the cosmic microwave background (CMB)—the oldest light in the Universe, released when (e) photons decoupled from matter and the Universe became (=) transparent to radiation |
| (703) | Until now, such a background of gravitational waves has (e) escaped detection. It is therefore very exciting that the BICEP2 (Background Imaging of Cosmic Extragalactic Polarization) collaboration has now reported in Physical Review Letters the observation of polarization modes of CMB radiation [1] that are consistent with (e&eb) such a background. |
| (704) | If confirmed, BICEP2's finding would provide (eb) the first direct probe of physical processes that occurred a mere 10−3610-36 seconds after (e) the big bang. |
| (705) | The result also offers a rare opportunity to directly test theoretical ideas, including inflation and grand unification (the theory that merges the electromagnetic, weak, and strong forces into a single force), at energies 13 orders of magnitude greater than (e) we can currently access at the Large Hadron Collider. |
| (706) | Inflation is based on (e) the possibility that a scalar field in the early Universe got momentarily stuck in (eb) a "false vacuum state"—a metastable condition in which the Universe wasn't at (eb) its global potential minimum. |
| (707) | The energy stored in this vacuum state could drive (eb) an exponential expansion of spacetime (similar to, but much faster than, the expansion due to dark energy seen today) before (e) a phase transition to the true vacuum state occurs, releasing (eb) this energy into (e&eb) particles and radiation. |
| (708) | This creates (eb) an initially hot state for (e) a subsequent big-bang-like expansion from (e) which our Universe subsequently evolved [2]. |
| (709) | If the Universe grew by at least a factor of 10261026, (eb) inflation can explain (eb) both the apparent flatness and isotropy of the Universe [2]: |
| (710) | The expansion would stretch (eb+ eb-dot) any initial curvature of (e) the Universe to near flatness, and smooth out (e&eb) any initial inhomogeneities, just as stretching (eb+ ebdot) a piece of fabric would do. |
| (711) | Isotropy is justified by the fact that the Universe, before inflation, was once much smaller and could have been contained in (eb) a causally connected region that was in thermal equilibrium, resulting in (eb) the approximately uniform, isotropic temperature distribution observed on (eb) large scales today. |

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| (712) Importantly, the theory can explain the (eb) origin of structure in the Universe: Microscopic quantum fluctuations associated with (e&eb) the scalar field driving inflation could have been "frozen in" during the exponential expansion, resulting, (eb) after inflation, in density fluctuations that might have seeded all observed large-scale structures in the Universe, including (e) all visible galaxies and clusters of galaxies. |
| (713) The signs of such primordial fluctuations have indeed been detected through (e&eb) their effect on (e&eb) the CMB, first by the Cosmic Background Explorer (COBE) [3] and later by other missions. |
| (714) These observational studies detected small temperature variations (anisotropies) in (eb) the CMB and found a spectrum compatible with (e&eb) that predicted by inflation. |
| (715) The magnitude of these fluctuations suggests (eb) an energy scale of inflation (loosely speaking, the energy stored throughout space in (eb) the scalar field potential that drove inflation) close to (e) 10161016 giga-electron-volts (GeV). |
| (716) But while consistent with inflation, such fluctuations are not (e) conclusive evidence: Their predicted shape and magnitude is known to vary with (e&eb) the details of the inflationary potential (whose form is not yet experimentally constrained). |
| (717) Hence many different observations could be accommodated, and any agreement between theory and observation is at best suggestive. There is also a possibility that other, noninflationary sources might contribute. However, gravitational waves—the propagating ripples in spacetime predicted by Einstein in 1916—might provide (eb) direct evidence of inflation. |
| (718) During inflation, gravitational fields would also fluctuate because of (e) quantum effects. |
| (719) These fluctuations would similarly get "frozen in" and re-emerge after inflation as (=) classical gravitational waves, whose amplitude depends only on (=) the energy stored in the field driving (eb) inflation, and not on other model-dependent details [4]. |
| (720) If the scale of inflation were indeed as high as a few times 10161016 GeV, this background could produce (eb) specific temperature anisotropies in the CMB [5]. |
| (721) These anisotropies have thus far not been detected by (e) the most recent searches by the WMAP and PLANCK satellites, and their absence was used to (e) put upper limits on the scale of inflation somewhat below 10161016 GeV. |
| (722) But a more sensitive CMB probe of (e) gravitational waves exists. |
| (723) Large-wavelength gravitational waves produced by (e) inflation would have alternatively compressed and expanded (e&eb) space, producing (eb) spatial variations in (eb) the intensity of radiation bombarding free electrons in (eb)the early Universe (see, e.g., Ref: [6]). |
| (724) As photons decoupled from (e) matter and CMB radiation was released, scattering from (e) the undulating sea of electrons would have created (eb) patterns in the polarization of CMB. |
| (725) Such polarization anisotropies are predicted in (eb) inflationary models to be (=) most visible when (e) looking at the sky on spatial scales comparable to (e&eb) or somewhat larger than the size of (e) the Universe's horizon at (eb) the time when the CMB was created. |
| (726) This corresponds to (e&eb) about 1 degree of the field of view of the sky. |
| (727) Over the past decade, various CMB polarization experiments have thus chased the observation of polarization anisotropies on (e&eb) these scales but, so far, also yielded (eb) merely upper limits on a possible gravitational background. |
| (728) All of this has now changed with (e&eb) the BICEP2 collaboration's [1] report of (e) a signal that has the predicted (eb) characteristics of such a background at precisely the angular scales where (e) the inflationary signal is expected (see Fig. 1, top deleted due to spatial restraints). |

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| (729) The BICEP2 Collaboration uses (e) a microwave telescope located at the South Pole, with (e&eb) an array of polarization-sensitive detectors pointing at (eb) a small field (a solid angle of about 11 percent of the full sky) with a resolution of approximately 0.50.5 degrees, thus sufficient to probe the relevant angular scales for the detection of (e) polarization anisotropies originating from (e) **inflation**. |
| (730) BICEP2's search is based on (e) two expected features of primordial gravitational waves. |
| (731) First, gravitational waves would produce (eb) a characteristic spatial polarization pattern on CMB—a consequence of (e) their quadrupole nature. |
| (732) One can decompose (e&eb) any spatial polarization pattern into (e&eb) two orthogonal components, a "curl-free" component and (e&eb) a component involving a twisting pattern. |
| (733) In analogy to electromagnetic modes, these polarization patterns are called "EE modes" and (e&eb) "BB modes," respectively (see Fig. 1). |
| (734) Scalar fluctuations in the density of (e) matter would produce (eb) EE modes, while (e) gravitational waves would instead generate (eb) a BB-mode pattern. |
| (735) The second feature is the aforementioned characteristic spatial scale of (e) expected CMB polarization signatures, including (e) wavelength scales comparable to (e&eb) or larger than the horizon scale at the time the CMB was created. |
| (736) Since inflation would produce (eb) both gravitational waves and (e&eb) scalar density perturbations, BICEP2, like all CMB-polarization experiments, quote their results in terms of (e&eb) the ratio rr of the amplitude of (e) tensor perturbations (associated with (e&eb) gravitational waves) to (e) the amplitude associated with (e&eb) scalar density fluctuations (reflected in (eb) the magnitude of CMB temperature fluctuations measured by (e) experiments like COBE). |
| (737) A nonzero rr implies (eb) the detection of a BB-mode pattern and primarily measures (eb) the strength of any possible gravitational wave signal which, if (e) it comes from inflation, is then determined by (e) its energy scale (e.g., see Ref. [7] for a review). |
| (738) BICEP2 reported a nonzero value of r=0.2(+0.07−0.05)r=0.2(+0.07-0.05), on scales ranging from (e) the horizon scale up to 4–5 times larger, with (e&eb) a claimed significance of better than 5σ5σ. |
| (739) If confirmed, a value of 0.2 corresponds to (e&eb) an energy density during (e&eb) inflation of ∼1.2×1016∼1.2×1016 GeV. |
| (740) The claimed signal is far larger than expected, partly because (e) the PLANCK satellite, despite not having yet reported its polarization measurements, has (e) previously put an upper limit on (eb) rr (<0.11<0.11). PLANCK's estimate is based on (e) the absence of a gravitational wave signature on (eb) the CMB temperature anisotropy on (eb) its largest scales (PLANCK observes the whole sky and therefore angular scales up to 360360 degrees). |
| (741) The apparent discrepancy between the PLANCK upper bound and (e&eb) the BICEP2 detection might be a cause for concern. |
| (742) However, the two bounds probe different spatial scales, and the BICEP collaboration has noted that, while tensor perturbations are expected to be (=) roughly scale independent, some scale dependence of (e) scalar density perturbations could relax (e&eb) the PLANCK constraint and bring (eb) the two results into agreement. |
| (743) Of greater concern, of course, is the possibility that the signal could be determined by (e) foreground contamination—any source along (e&eb) the line of sight of the detector that could distort or contaminate (e&eb) the primordial CMB signal. |
| (744) BICEP2 tested for (e) a wide variety of possible foregrounds, claiming (eb) they cannot be responsible for (e) the observed signal. |

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| (745) One obvious candidate arises from (e) gravitational lensing: The intervening mass of (e) astrophysical bodies could distort (e&eb) the spatial pattern of polarization in a way that that could transform (e&eb) an EE-mode signal into (e&eb) a BB-mode signal. |
| (746) The BICEP collaboration has analyzed this possibility extensively. Based on (e) measurements of (e) CMB temperature (and hence density perturbations) and on (e&eb) models of the mass distribution, they argue that such effects would only be (=) significant on smaller angular scales than those contributing to their measurements of rr. |
| (747) But since BICEP2's first announcement on March 17, 2014, a number of scientists have raised another foreground concern due to (e) the possible impact of (e&eb) our Galaxy's dust. |
| (748) BICEP2 modeled dust contamination scenarios based on (e) a variety of theoretical models, using data from (e) its predecessor BICEP1 and preliminary dust estimates from (e) the PLANCK satellite. |
| (749) They concluded that dust could not (e) reproduce the magnitude of the observed signal. |
| (750) However, PLANCK has in the meanwhile released (eb) more complete measurements of variation in polarized dust foregrounds in (eb) many regions of the Galaxy. |
| (751) Their analysis increases the uncertainty associated with (e&eb) the dust contribution to (e) the BICEP2 result and also accommodates (e) a potentially larger polarized dust contribution. |
| (752) Some researchers argued that dust polarization at a level 2–3 times higher than that assumed by BICEP2 is possible and could therefore plausibly be responsible for (e) the entire signal (e.g., see Ref: [8]). |
| (753) While the BICEP2 collaboration continues to maintain that this is extremely unlikely based on (e) their estimates, many other researchers are more skeptical. |
| (754) In any case, a variety of upcoming measurements, from PLANCK and other detectors are (=) expected to resolve this important ambiguity within the coming year. |
| (755) Even if the BICEP2 signal does arise from (e) primordial gravitational waves, various consistency checks will be required (e) before it can be definitively argued that the signal arises from (e) inflation. |
| (756) Inflation is certainly the best-motivated possibility, and the spectral shape of the signal, at least for (e) the five largest-scale modes observed by BICEP2, agrees extremely well with (=) inflationary predictions (see Fig. 1, bottom). |
| (757) Several results have already appeared that strongly constrain (e) other possible sources, ruling out (e) for example, the possibility of gravitational waves generated by (e) other possible phase transitions after inflation [9] |
| (758) . This conclusion is reinforced by more careful analyses of the scale dependence of (e) the predicted signal (e.g., Ref. [10]). |
| (759) But it will nevertheless be important to assure that all possible alternative sources of (e) gravitational waves are ruled out. |
| (760) If it can be established that the BICEP2 represents (eb) a gravitational wave signal from (e) inflation, there is one additional exciting implication. |
| (761) Standard calculations suggest (eb) that the gravitational wave signal from inflation is due to (e) quantum fluctuations in the gravitational field, and this appears to require (e) the existence of gravitons—the hypothetical elementary particles that mediate (e&eb) the force of gravity. |
| (762) Research by Frank Wilczek and myself, based on (e) a dimensional analysis argument [11], suggests that,(eb) independently of the method used to calculate the gravitational wave spectrum, a quantum gravitational origin is (=) required. |
| (763) If so, the BICEP2 result then implies (eb) that gravity is ultimately a quantum theory, a result of fundamental importance for physics. |

(764) It is hard to imagine a more important or wider-ranging cosmological observation. If confirmed by upcoming experiments, the BICEP2 results will push the frontiers of empirical physics forward in an unprecedented way, offering (eb) the first direct observation of gravitational waves and of quantum gravity effects as well as a new probe of the earliest moments in the Universe's life and of physics at grand-unified energy scales. This will thus constrain both our understanding of fundamental interactions and the evolution of our Universe as a whole.

**Detection of B-Mode Polarization at Degree Angular Scales by BICEP2 P. A. R. Ade et al. (BICEP2 Collaboration) Phys. Rev. Lett 112, 241101 – Published 19 June 2014**

(765) Authors report results from the BICEP2 experiment, a cosmic microwave background (CMB) polarimeter specifically designed to search for (e) the signal of inflationary gravitational waves in (eb) the B-mode power spectrum around $\ell \sim 80$.

(766) The telescope comprised a 26 cm aperture all-cold refracting optical system equipped with (e&eb) a focal plane of 512 antenna coupled transition edge sensor 150 GHz bolometers each with temperature sensitivity of $\approx 300\ \mu K_{CMB}\sqrt{s}$. BICEP2 observed from (e) the South Pole for three seasons from 2010 to 2012.

(767) A low-foreground region of sky with (e&eb) an effective area of 380 square deg was observed to (e) a depth of 87 nK deg in Stokes Q and U.

(768) In this paper authors describe the observations, data reduction, maps, simulations, and results. They find an excess of B-mode power over the base lensed-ΛCDM expectation in the range $30<\ell<150$, inconsistent with (e&eb) the null hypothesis at (eb) a significance of $>5\sigma$.

(769) Through jackknife tests and simulations based on (e) detailed calibration measurements authors show that systematic contamination is (=) much smaller than the observed excess.

(770) Cross correlating against WMAP 23 GHz maps we find that Galactic synchrotron makes a negligible contribution to the observed signal. We also examine a number of available models of polarized dust emission and find that at their default parameter values they predict power $\sim(5–10)\times$ smaller than the observed excess signal (with no significant cross-correlation with our maps). However, these models are not sufficiently constrained by external public data to exclude the possibility of dust emission bright enough to explain the entire excess signal. Cross correlating BICEP2 against 100 GHz maps from the BICEP1 experiment, the excess signal is confirmed with $3\sigma$ significance and its spectral index is found to be consistent with that of the CMB, disfavoring dust at $1.7\sigma$. The observed B-mode power spectrum is well fit by a lensed-ΛCDM+tensor theoretical model with tensor-to-scalar ratio $r=0.20+0.07-0.05$, with $r=0$ disfavored at $7.0\sigma$. Accounting for the contribution of foreground, dust will shift this value downward by an amount which will be better constrained with upcoming data sets. Received 4 April 2014 DOI:http://dx.doi.org/10.1103/PhysRevLett.112.241101 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**(771) Physicist claims to have observed quantum effects of Hawking radiation in the lab for the first time August 16, 2016 by Bob Yirka report**

(772) This computer-simulated image shows a supermassive black hole at the core of a galaxy. The black region in the center represents the black hole's event horizon, where no light can escape the massive object's gravitational grip. The black …more

(773) (Phys.org)—Jeff Steinhauer, a physicist at the Israel Institute of Technology, has published a paper in the journal Nature Physics describing experiments in which he attempted to create a virtual black hole in the lab in order to prove that Stephen Hawking's theory of radiation emanating from black holes is correct —though his experiments are based on sound, rather than light. In his paper, he claims to have observed the quantum effects of Hawking radiation in his lab as part of a virtual black hole—which, if proven to be true, will be the first time it has ever been achieved.

(774) For many years, scientists believed that nothing could ever escape from a black hole. But in 1974, Stephen Hawking published a paper suggesting that something could—particles that are now called Hawking radiation. His idea was that if a particle (and its antimatter mate) appeared spontaneously at the edge of a black hole, one of the pair might be pulled into the black hole while the other escaped, taking some of the energy from the black hole with it—which would explain why black holes grow smaller and eventually disappear. Because such emissions are so feeble, no one has been able to measure Hawking radiation, so researchers have instead tried to build virtual black holes in labs to test the theory. One type of virtual black hole was proposed back in 1981 by Bill Unruh with the University of British Columbia—he suggested that an analogue might be created using water instead of light. He imagined a phonon existing at the edge of a waterfall—as the water speeds up, it begins to move faster than the speed of sound, causing it to be trapped. But if the phonon had an entangled mate that eluded the fall by moving away before getting caught up, it could escape. In this new effort, Steinhauer has built a device based on that idea and in so doing, claims he has observed an analogue of Hawking radiation.

(775) The experiment consisted of creating an entangled pair of phonons sitting inside a bit of liquid that had been forced (via laser) to move very fast and then observing the action as one of the pair was pulled away as the liquid began to move faster than the speed of sound, while the other escaped—the fluid was a Bose-Einstein condensate of rubidium-87 atoms. After repeating the experiment 4,600 times Steinhauer became convinced that the particles were entangled, a necessity for a Hawking radiation analogue. His findings do not prove Hawking's theory to be true, of course, but they do appear to add a degree of credence that other researchers have thus far not been able to achieve. More information: Jeff Steinhauer. Observation of quantum Hawking radiation and its entanglement in an analogue black hole, Nature Physics (2016). DOI: 10.1038/nphys3863

**(776) Abstract**

We observe spontaneous Hawking radiation, stimulated by quantum vacuum fluctuations, emanating from an analogue black hole in an atomic Bose–Einstein condensate. Correlations are observed between the Hawking particles outside the black hole and the partner particles inside. These correlations indicate an approximately thermal distribution of Hawking radiation. We find that the high-energy pairs are entangled, while the low-energy pairs are not, within the reasonable assumption that excitations with different frequencies are not correlated. The entanglement verifies the quantum nature of the Hawking radiation. The results are consistent with a driven oscillation experiment and a numerical simulation. Journal reference: Nature Physics search and more info website

Read more at: http://phys.org/news/2016-08-physicist-quantum-effects-hawking-lab.html#jCp

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as

the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t - T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

These equations can be written for

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(4) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.

(5) However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].

(6) Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.

Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

**TIME DELAY**

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period,

which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types system takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t - T_7)F_7(R_7(t - T_7) - N_7(t)M_7((R_7(t) \qquad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7}+\frac{N_7 dG_7}{dR_7}-\frac{dB_7}{dR_7}>\frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7}+\frac{dM_7}{dR_7}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also normalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_7$) and the other derived from the non-delayed or "immediate" processes($I_7$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_7=\frac{N_7 G_7}{2M_7}\frac{d(7nF_7)}{dR_7}(\geq 0) \quad (7a)$$

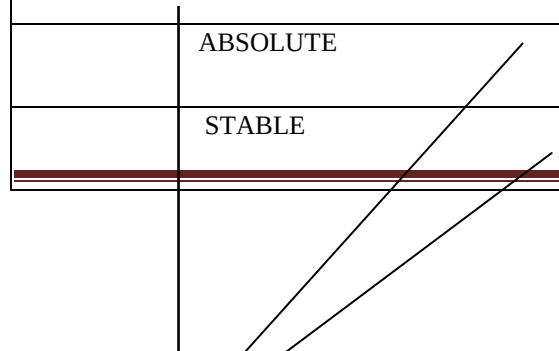
$$I_7=-\frac{N_7 G_7}{2M_7}\frac{d(7nM_7}{dR_7}(\geq 0) \quad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7>L_7-I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

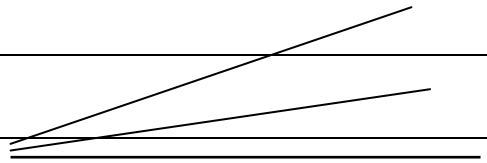
$$P_7>L_7-I_7, \quad \text{when} \quad P_7^2/4>I \quad (8a)$$

$$P_7>L_7-I_7+\frac{(I_7-P_7^2}{(L_7-I)} \quad \text{when} \quad P_7^2/4<I. \quad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,





**DELAYED PREDATOR CREDIT TYPES EFFECT WITH ITS WIDE RANGING RAMIFICATION**

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

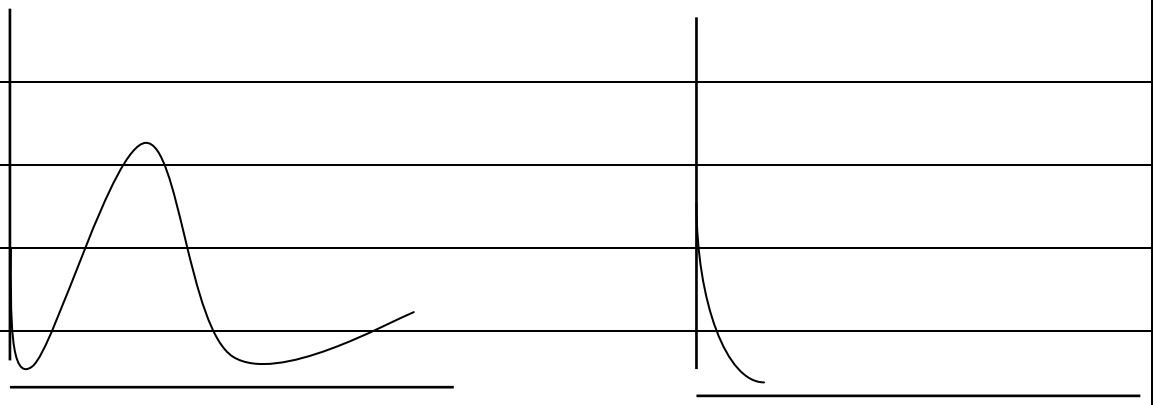
It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay

and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



**TIME DELAY**

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

To scores of problems this relationship is important. Of them some of the examples are as follows:

**New look at the galaxy-shaping power of black holes**

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| (129) | Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving. |
| (130) | "These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science. |
| (131) | "Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy." |
| (132) | This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return. |
| (133) | The remaining matter gets absorbed into the black hole, adding to (e) its mass. |
| (134) | The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne. |
| (135) | The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution. |
| (136) | Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form. |
| (137) | "Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara. |
| (138) | "It creates (eb) these bubbles, and the bubbles keep the plasma hot. |
| (139) | That's what prevents galaxies from (e) becoming even bigger than they are now." |
| (140) | Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space. |
| (141) | In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene. |
| (142) | Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve. |
| (143) | "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara. |
| (144) | "It's like nature's thermostat, which keeps these galaxies from (e) growing. |
| (145) | If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars." |

**A new look at the galaxy-shaping power of black holes**

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| (146) | An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the |
| (147) | McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said. |

(148) Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.

**Doomed satellite missions**

(149) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.

(150) At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.

(151) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.

(152) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.

(153) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.

(154) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people

(155) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."

(156) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.

(157) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(158) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.

(159) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.

(160) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.

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| (161) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses. |
| (162) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well. |
| (163) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma. |
| (164) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6. |
| (165) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus. |
| (166) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (167) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (168) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (169) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (170) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (171) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (172) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (173) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (174) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (175) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (176) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (177) We know dark matter exists only from (e) its gravitational effects. |
| (178) It holds together galaxies and (e&eb) galaxy clusters. |

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| (179) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (180) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |
| (181) They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (182) X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (183) They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (184) This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (185) An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |
| (186) Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (187) "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (188) The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (189) Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (190) Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (191) "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (192) Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (193) "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New |

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| York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (194) Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories. |
| (195) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen. |
| (196) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer. |
| (197) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants. |
| (198) These were treated as (=) "aware" and "unaware" trials, respectively. |
| (199) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it. |
| (200) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another. |
| (201) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion. |
| (202) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness. |
| (203) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said. |
| (204) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness." |
| (205) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner. |
| (206) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain. |
| (207) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said. |
| (208) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

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| (209) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter. |
| (210) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined. |
| (211) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions. |
| (212) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda. |
| (213) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general. |
| (214) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole. |
| (215) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes. |
| (216) Noda and his colleagues instead studied the change in emissions over (e&eb) time. |
| (217) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection. |
| (218) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1). |
| (219) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole. |
| (220) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved. |
| (221) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp |

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

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| (222) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking “palm tree” organization. |
| (223) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society |
| (224) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible? |
| (225) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem. |
| (226) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other. |
| (227) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible. |
| (228) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges. |
| (229) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century. |
| (230) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching. |
| (231) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom. |
| (232) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth. |

**Spheres on trees**

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| (233) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration. |
| (234) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement. |
| (235) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum. |
| (236) By visualizing the problem in this way, the researchers noticed that these particular graphs don’t (e) have lots of branches extending from lots of (e) other branches. |

| |
|---|
| (237) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure. |
| (238) The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network. |
| (239) As a result, it provides (eb) an important clue for finding the global minimum. |
| (240) It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum. |
| (241) This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large. |

**Thompson's problem 2**

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| (242) By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society |
| (243) The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima. |
| **(244)** The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node. |

**Implications for Thomson's and Smale's problems**

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| (245) Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before. |
| (246) "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around." |
| (247) As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem. |
| (248) "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view." |
| (249) In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges. |
| (250) From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear. |
| (251) "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK. |
| (252) "For larger numbers, there will likely be (=) multiple funnels. |
| (253) Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum." |

(254) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.

(255) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.

(256) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t-T_7)F_7\big(R_7(t-T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t-T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) and the effective prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this

rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Look at the resemblance of String Theory and Abhinava's earlier theory and how the model could be applied: The world of manifestation and differentiated perceptions (Yikalpa) thus extends from one Centre to (eb) the next. **What is meant again is the realisation or the evolution of individual consciousness or general ledger to be (=)the cosmic general ledger is the sole criterion for the Satchidananda state where all dualities cease** Although it is never in fact divorced from (e) the subject who resides there, the ignorant fail (e) to grasp this fact and so, cut off from (e) the Centre, the world of objectivity becomes (=) for them the sphere of Maya. Bhagavatotpala quotes the Light of Consciousness amvitprakdsa): This ever pure experience (suddhanubhava) is variegated by (e) each form [revealed within it] Even so it remains unstained (nirmala) when moving to another. Just as a cloth which is naturally white, once dyed, cannot (e) change colour without [first] becoming white again, similarly the pure power of awareness, (citi) once coloured by (e) form, is pure [again] at the Centre where that form is (=) abandoned and from whence it proceeds to (e&eb) another. In his Essence of Vibration (Spandasarpdoha), K§emaraja explains that the rise and fall of every individual perception in the field of awareness is (=) a specific pulsation of consciousness. This is exactly the individual general ledger transaction lifted to the transaction in cosmic general ledger. During The Initial Instant Of Perception, T Consciousness Is Manifestly Apparent And The Yogi, Participating In Its Plenitude, Observes The Outer World Without Being Attached To Any Particular Divine Body And Sacred Circle Of The Senses Singling It Out From Any Other, Like A Man Who Observes A City From A High Mountain Peak.

He Sees The Outer World Reflected Within His Consciousness Free Of Thought-Constructs And So 'Stamps* The Outer On The Inner While Absorbing The Object And Means Of Knowledge In The Pure Subject Which Grasps Them As The Expansion Of His Own Nature. K§emaraja Says: By Penetrating Into Bhairavimudra, The Yogi Observes The Vast Totality Of Beings Rising From, And Dissolving Into, The Sky [Of Conscious- Ness], Like A Series Of Reflections Appearing And Disappearing Inside A Mirror. Through (e&eb) The Practice Of Bhairavimudra The Yogi Realises That He Is (=) The Substratum Consciousness (Adhisfhatr) Which Both Underlies (e&eb) And Is (=) The Essence Of All Things. He Discovers That Phenomena Have (e) No Independent Existence Apart From (e) Him And So Are, In This Sense, Void. At The Same Time, He Realises That Because (e) All Things Are Consciousness, They Are (=) Far From Unreal. He Views The Outer World Yet Sees (e) It Not. Beyond Both Voidness And Non-Voidness He Penetrates Into (e&eb) The Supreme Abode (Pararn Padam) Of Siva's Consciousness. [The Powers Of The Senses] Endowed With (e&eb) The Attributes Of The Great Union [Between Subject And (e&eb) Object] Whose Form Is The Wakening Of Man's Spiritual Potential (Kuntfalinf), Fill (e&eb) [With Consciousness] The Outer Clatter Of Diversity (Bhedatfambara) Born Of (e) Its Intense Power And Are Then Established In (eb) The Unobscured Abode Of The Void Of Consciousness To Shine [There] Eternally. Thus Residing Beyond (e) Being And Non-Being, The Sole Protector Of The Unity Which Is Tranquil And Expanding [Consciousness], Whose Glory Is All-Embracing And Form Unobscured, Is Called Bhairavimudra. Through (e&eb) The Practice Of Bhairavimudra, The Yogi Unites The Universal Vibration Of T Consciousness With (e&eb) The Individual Pulsation Of Objectivised 'This' Consciousness. The Two Aspects Of Consciousness Are Now In (eb) A State Of Equilibrium Like (=) The Two Pans Of An Evenly Weighed Balance And The Yogi Experiences (e&eb) The Pure Knowledge (Suddhavidya) That: 'I Alone Am (=)All Things'. Thus Becoming (=) The Master Of The Wheel Of Energies He Is (=) Free, Like Siva, To Create And Destroy. When [The Yogi] Is Well Established, Without (e) Wavering, Solely In The Integral Egoity Of His Authentic Nature, The Spanda Principle, And Is

Absorbed In Contemplation (Samavisfa), He Becomes (=) One With It The Doctrine Of Vibration (Tanmaya). Then . . . Dissolving And Creating The Universe By Means Of (er) His Introverted And Extroverted Absorption, He Destroys And Creates (e&eb) All Things Out Of Sarikara, His Innate Nature. [Thus] He Assumes The State Of The Universal Experiencer And Having Absorbed All That Is To Be Experienced From [The Grossest Level] — Earth — To (The Subtlest) — Siva — He Reaches (eb) The State Of The Supreme Subject By (e) Progressively Recognising [His Identity With Him]. Thus, Introverted And Extroverted Absorption Both Lead To (eb) The Recognition Of The Pulsation (Spanda) Of (e) One's Own Consciousness. This Is Where (e) Individual General Ledger Becomes (=) Cosmic General Ledger. At The Level Of Consciousness Corresponding To (e&eb) Siva's Basic State (Saryibhavdvastha), The Alternation From Inner To (e&eb) Outer Is Instantaneously Resolved Into (e&eb) The Vibration Of His Nature. When The Yogi Finally Comes To Be (=) Constantly Aware Of This Reality, His Enlightenment Is (=) Full And Perfect. Freed Of All Means (Anupaya) And Delighting In The Power Of His Bliss (Anandasakti), He Knows And Does Whatever He Pleases. The Yogi Seeking Self-Realisation Must Acquire (e) Mastery Over This Movement. Ksemaraja Stresses That The Doctrine Of Vibration Teaches That Liberation Can Only Be Achieved By (e) First Withdrawing All Sense Activity In Introverted Contemplation (Nimilanasamadhi) To (e) Then Experience The 'Great Expansion' (Mahdvikdsd) Of (e) Consciousness While (e) Recognising This To Be (=) A Spontaneous Process Within (eb) It. This Is Done Through (e&eb) The Practice Of Kramamudra. A Passage From The Now Lost Kramasiitra Explains: Although The Adept's Attention [May Be] Outwardly Directed, (eb) He Enjoys Contemplative Absorption Through (e&eb) The Introverted Aspect Of (e) Kramamudra. Initially He Turns Inward From (e)The Outside World And [Then] From Within [Himself] He Exits Into The Outer World Under The Influence Of His Absorption. There Is No Necessity Of Any Glorification Or Mortification Here. It Is Just Like Joining A Job And Retiring And Realising All That Happened Is Just By His Own Effort And Nothing Else. Thus The Sequence (Krama) In This Attitude (Mudra) [Ranges Through] Both Inner And Outer. The Yogi Must Pervade

The Surface Level Of Awareness (Yyutthana) With The Same Bliss He Experiences Plunging Into The Depths Of Con- Templative Absorption (Samadhi). Note Here That Samadhi Is A Death Like State And Still Being Matter And Information Stored You Experience Consciousness. Submerging Himself And Emerging Repeatedly From Samadhi, He Eventually Recognises That The Unity Of Consciousness Pervades Both States: The Best Of Yogi's, Who Has Achieved A State Of Complete Absorption Even When Risen From Meditation, [Inwardly] Vibrating Like A Drunkard In Blissful Inebriation From The After-Effects Of The Nectar Of Con- Templation, Sees All Things Dissolving In The Sky Of Consciousness Like A Cloud In The Autumn Sky. He Plunges Repeatedly within Himself Divine Body and Sacred Circle of the Senses and Becomes (=) Aware Of His Identity With (e&eb) Consciousness By (e) The Practice Of Introverted Contemplation. Thus Even When He Is Said To Have Risen From (e) Absorption, He Is (=) One With [His] Experience Of It. Ksemaraja Goes On To Explain That This Practice Is Called 'Mudra' Because (e) It Both Fills The Adept With (e&eb) Bliss (Mud) And Is Itself The Bliss Of (e) Consciousness. Moreover, It Dissolves (e) Away (Dra) All Bondage And 'Stamps' The Universe Of Experience With (e&eb)The Seal (Mudra) Of The Fourth State (Turiya) Of Enlightened Consciousness Beyond, And Including, The Three States Of Waking, Dreaming And Deep Sleep. It Is Called 'Krama' Because (e) It Is The Root Source Of (e) All Emanation And All Other Conscious Processes Which Succeed (e&eb) One Another In Ordered Sequence (Krama) And Is, (=)At The Same Time, Their Successive (Krama) Appearance As Well. By The Practice Of Kramamudra The Opposites Fuse And Siva And Sakti Unite. They Yogi Comes To Experience The Simultaneous Pervasion Of All The Lower, Grosser Categories Of Existence By The Higher And The Presence Of The Lower In The Higher. Commencing His Practice In A Low Form Of Bhairavimudra, The Yogi Conjoins The Outer With The Inner; Then, In Kramamudra, He Fills Both The Outer With (e&eb) The Inner And (e&eb) The Inner With (e&eb)The Outer. When He Achieves (eb) Perfection In This Two-Fold Movement, He Attains To (e) The Highest Form Of Bhairavimudra In Which The Two Merge Completely In (eb)The Experience Of (e) The Absolute (Anuttara), Free

Of (e) All Differentiation And (e&eb) Polarities. If He Fails To Maintain Awareness Of This State, He Again Falls Into (eb) Kramamudra Until He Has Finally Completely Merged All The Highest States In (e&eb) The Lower And The Lower In (e&eb) The Higher. He Then No Longer Needs To (e) Resort To (e) Any Means (Anupaya) To (e) Achieve Liberation. All He Says or Does Anything He Perceives or Thinks, Instantly Occasions In (eb) Him The Highest Level Of Consciousness. Thus The Fruit Of Bhairavimudra Is (*=)The Wonder (Camatkara) Or Amazement (Vismaya) That Overcomes (e)The Yogi When (e) He Reaches The Plane Of (e) Union (Yogabhumika) Y Xsl Where (e) All Opposites Merge In (e&eb)The Radiance Of The Great Light Of Consciousness. The Stanzas On Vibration Teach: How Can One Who, As If Astonished, Beholds (e) His Own Nature As (=) That Which Sustains [The Existence Of Everything] Be Subject To (e&eb) This Painful Round Of Transmigration (Kusrti)? The Yogi, Recognising His True Nature To Be (=)The Supreme Subject, Is Astonished To (e) Suddenly Discover That The Individual He Thought He Was, Caught Up (e&eb) In The Trammels Of Thought And Living In A World Enmeshed In (e&eb) The Web Of Time And Space, Does Not Really Exist At All. He Experiences A 'Turning About' (Paravrtti) In (eb)The Deepest Seat Of The Doctrine Of Vibration Consciousness As He Penetrates His True Nature. The Sudden Eruption Of This Intuition (Pratibha) Arouses In Him A Cry Of Amazement As He Transcends All Thought-Constructs And, Perfectly Absorbed In His Own Nature, Is Liberated. The Path To Liberation Essentially, Spanda Doctrine Is Concerned With Two Matters. The First Is To Impart To Those Who Are Fit To Receive The Teachings A Deeper Understanding Of The Ultimate Goal Of Life (Upeya). When We Have Understood What Truly Benefits Us And Is Worth Attaining And What, On The Contrary, Is Of No Real Value But Stands In The Way Of This Attainment, We Can Begin To Make Progress Towards Our Goal. This Is Spanda Doctrine's Second Concern, Namely, To Show The Way In Which We Can Develop Spiritually Through Siva's Grace And The Right Application Of The Means To Realisation That It Teaches. When Both These Aspects Of The Teaching Have Been Correctly Understood And Applied, The Spanda Yogi Achieves A Clear And Permanent

Realisation Of His Goal And Is Liberated, Thus Fulfilling The Ultimate Aim Of The Teaching. The Doctrine Of Vibration Is Not Meant For The Spiritually Dull. It Is Not For The Worldly Whose Consciousness, Clouded By Ignorance, Is As If Dreaming, Even During The Waking State Of Daily Life, The Dream Of Its Own Thought-Constructs. The Teachings Are Meant For Those Who Are Awake (Prabuddha), Those Who, Full Of Faith And Reverence, Are Always Alert And Intent On Discerning The True Nature Of Ultimate Reality. This Reality Is Understood In Three Basic Ways. The First Is Purely Transcendental. The Stanzas Choose This Aspect As The One Which Formally Defines It Most Specifically. Ultimate Reality Transcends All The Opposites, Including Subject And Object. This Does Not Mean, However, That It Is An Unconscious Void, A Mere Absence Of All Existence. In Fact, This Negative Characterization Of Reality (Which Includes Also A Denial Of All That Is Unconscious) Implies A Positive Immanence In Which The Opposites Are United In The Oneness Of Pure Consciousness That Is Equally Siva And Spanda, His Universal Activity. These Two Seemingly Contrasting Aspects Are Reconciled In The Third, Namely, Reality Understood As The Essential Nature Of All Things. Although Universal And Everywhere The Same, It Is Understood To Be The Essential And Specific Nature Of Each Existent As Its 'Own Nature' (Svabhava).

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such

a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

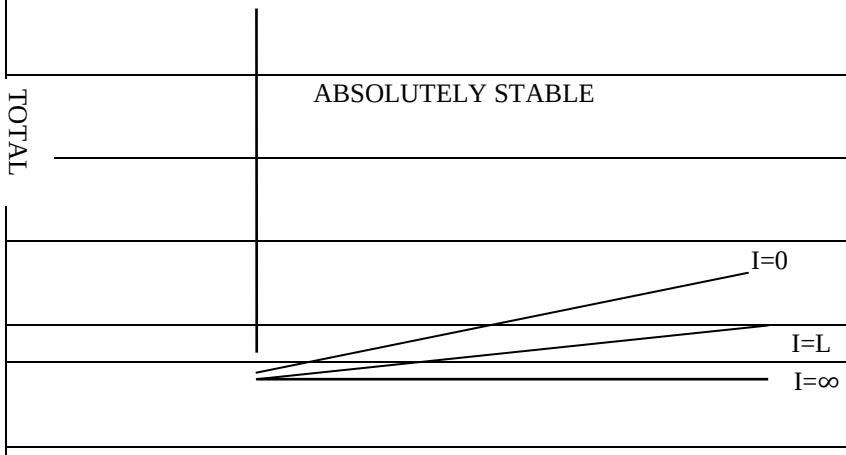
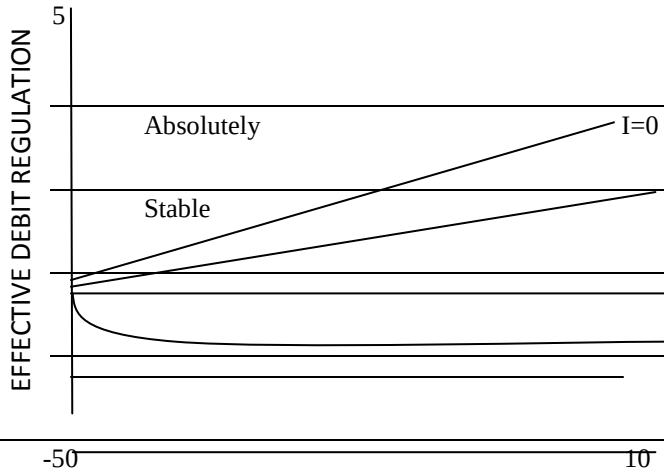



FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

This equation is of paramount importance in the following cases: According to the principles of geometric quantum mechanics, the physical characteristics of a given quantum system can be represented by geometrical features that are preferentially identified in this complex manifold. Here we construct a number of examples of such features as they arise in the state spaces for spin, spin 1, spin and spin 2 systems, and for pairs of spin systems. A study is then undertaken on the geometry of entangled states. A locally invariant measure is assigned to the degree of entanglement of a given state for a general multi-particle system, and the properties of this measure are analysed for the entangled states of a pair of spin particles. With the specification of a quantum Hamiltonian, the resulting Schrödinger trajectories induce an isometry of the Fubini–Study manifold, and hence also an isometry of each of the energy surfaces generated by level values of the expectation of the Hamiltonian. For a generic quantum evolution,

the corresponding Killing trajectory is quasiergodic on a toroidal subspace of the energy surface through the initial state. When a dynamical trajectory is lifted orthogonally to Hilbert space, it induces a geometric phase shift on the wave function. The uncertainty of an observable in a given state is the length of the gradient vector of the level surface of the expectation of the observable in that state; a fact that allows us to calculate higher orders corrections to the Heisenberg relations. A general mixed state is determined by a probability density function on the state space, for which the associated first moment is the density matrix. The advantage of a general state is in its applicability in various attempts to go beyond the standard quantum theory, some of which admit a natural phase-space characterization. **Journal of Geometry and Physics Volume 38, Issue 1, April 2001, Pages 19–53 Geometric quantum mechanics Dorje C. Brody Lane P. Hughston DOI: 10.1016/S0393-0440(00)00052-8** Expansion (unmesa) of [consciousness] or the creative intuition {pratibha) is (=) [experienced] in the interval which divides two [moments] of differentiated perception (vikalpa) It is here that they arise and disappear. The Sastras and Agamas proclaim with reasoned argument that it is (=) free of thought-constructs {Nirvikalpa) and precedes (e) Path to Liberation all mental representations of any object. None can deny that a gap exists between perceptions insofar as two moments of thought are (=) invariably divided. This [gap] is (=) the undifferentiated unity of all the countless manifestations. Similarly, in the outer more objective sphere, where change consists (e) of the alterations in the configurations of manifest appearances (Abhasa\ the transition from one to (e&eb) another corresponds to (e&eb) a phase of pure luminosity that marks (eb) the beginning of one form and the end of another

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^{\infty} \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7 \left( U_7 - V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7 \right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$then given the constraints of $f(T_7), (A1.5) can\ never\ be\ satisfied\ Co$nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7 \lambda_7 e^{-T_7 \lambda_7} + [(U_7 - V_7) M_7 + Y_7] - [((U_7 -$$

V 7M 7–]**e**–7 7λ 7=0

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$\boldsymbol{Z}_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7) \cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7) \sin(Z_7 T_7) = 0$$

The warrant of Truth is ever elusive when we deconstruct the foundational justifications of our convenient notions about the way the world works. It is easy to confuse what is actually the creation of beliefs with the "discovery of Truth," a common goal of science and theology. Science offers no ontological Scientific Picture of the World but substitutes a number of applied theories as meta-theory. It's close, powerful and useful: it 'works' and allows use to do work. But the still mysterious "true" metatheory may reveal the basis of consciousness. But who knows when we will lift that veil? Permeating the living reality of our culture are certain contagious notions, fads and trends that have the ability to influence the way we think about the nature of Reality and ourselves. Some of them are toxic and can consume you; others are just whims of pop culture infectious and benign as the common cold. **Jung** described the concept of psychic contagion by certain archetypal forces inherent in the human psyche, which manifest in our spiritual lives and belief systems. He spoke of both conscious and unconscious contamination. Notions like this range from simple superstitions to scientific concepts, to urban myths. Notions sweep through our culture and insinuate themselves within its fabric, as fads, whether they are "real" or not, they can be influential. They mood alter us, make us feel we belong. An analysis of these notions is useful in distinguishing a common human phenomenon from any potential "alien influence" which may or may not be exerted on us from an unknown source. **Strange Attractors: Transference, Holography, and an Archetype Burke, J. (2003).**

**Strange Attractors: Transference, Holography, and an Archetype (Doctoral dissertation, Pacifica Graduate Institute, 2003).** In consideration to the Nobel Proze for Chemistry this year(2014), a stability nalysis of the proble in to single molexues would be a contemporaneous question that has to be taken cognisace of and given importance to . Determination of molecular structure by geometry optimization became routine only after efficient methods for calculating the first derivatives of the energy with respect to all atomic coordinates became available. Evaluation of the related second derivatives allows the prediction of vibrational frequencies if harmonic motion is estimated. More importantly, it allows for the characterization of stationary points. The frequencies are related to the eigenvalues of the Hessian matrix, which contains second derivatives. If the eigenvalues are all positive, then the frequencies are all real and the stationary point is a local minimum. If one eigenvalue is negative (i.e., an imaginary frequency), then the stationary point is a transition structure. If more than one eigenvalue is negative, then the stationary point is a more complex one, and is usually of little interest. When one of these is found, it is necessary to move the search away from it if the experimenter is looking solely for local minima and transition structures. The total energy is determined by approximate solutions of the time-dependent Schrödinger equation, usually with no relativistic terms included, and by making use of the Born–Oppenheimer approximation, which allows for the separation of electronic and nuclear motions, thereby simplifying the Schrödinger equation. This leads to the evaluation of the total energy as a sum of the electronic energy at fixed nuclei positions and the repulsion energy of the nuclei. Notable exceptions are certain approaches called direct quantum chemistry, which treat electrons and nuclei on a common footing. Density functional methods and semi-empirical methods are variants on the major theme. **Orlando Alvarez** discusses Polyakov's quantization of the string in the presence of a boundary allowing for an arbitrary topology for the world sheet. In addition to the dynamical conformal factor discovered by Polyakov, there are a finite

number of new degrees of freedom if the surface is more complicated than a sphere or a disc. The quantization of the Liouville theory in an arbitrary topology is also discussed. A one-loop calculation shows that the model is renormalizable if one performs a mass renormalization and an additive field renormalization. The renormalization group equations have a perturbative infrared unstable fixed point in all topologies. Copyright © 1983 Published by Elsevier B.V. **Nuclear Physics B Volume 216, Issue 1, 25 April 1983, Pages 125–184 Theory of strings with boundaries: Fluctuations, topology and quantum geometry Orlando Alvarez DOI: 10.1016/0550-3213(83)90490-X** Quantum geometry (the modern loop quantum gravity involving graphs and spin-networks instead of the loops) provides microscopic degrees of freedom that account for black-hole entropy. However, the procedure for state counting used in the literature contains an error and the number of the relevant horizon states is underestimated. In our paper a correct method of counting is presented by **Marcin Domagala and Jerzy Lewandowski.** Results lead to a revision of the literature of the subject. It turns out that the contribution of spins greater than 1/2 to the entropy is not negligible. Hence, the value of the Barbero–Immirzi parameter involved in the spectra of all the geometric and physical operators in this theory is different than previously derived. Also, the conjectured relation between quantum geometry and the black-hole quasi-normal modes should be understood again. **Marcin Domagala and Jerzy Lewandowski 2004 Class. Quantum Grav 21 5233 doi:10.1088/0264-9381/21/22/014 Black-hole entropy from quantum geometry** Quantum geometry predicts that a universe evolves through an inflationary phase at small volume before exiting gracefully into a standard Friedmann phase. This does not require the introduction of additional matter fields with ad hoc potentials; rather, it occurs because of a quantum gravity modification of the kinetic part of ordinary matter Hamiltonians. Authors draw the cognizance of application of the same mechanism that can explain why the present day cosmological acceleration is so tiny. DOI: http://dx.doi.org/10.1103/PhysRevLett.89.261301

**Inflation from Quantum Geometry Phys. Rev. Lett 89, 261301 – Published 12 December 2002 Martin Bojowald** The loop quantum cosmology of the closed isotropic model is studied with special emphasis on a comparison with traditional results obtained in the Wheeler-DeWitt approach by **Martin Bojowald** .Investigation includes the relation of the dynamical initial conditions to boundary conditions such as the no-boundary or the tunneling proposal and a discussion of inflation from quantum cosmology. **DOI: http://dx.doi.org/10.1103/PhysRevD.67.124023 Loop quantum cosmology, boundary proposals, and inflation Phys. Rev. D 67, 124023 – Published 23 June 2003 Martin Bojowald and Kevin Vandersloot** In loop quantum cosmology, the universe avoids a big bang singularity and undergoes an early and short super-inflation phase. During super-inflation, non-perturbative quantum corrections to the dynamics drive an inflaton field up its potential hill, thus setting the initial conditions for standard inflation. **Shinji Tsujikawa et al** show that this effect can raise the inflaton high enough to achieve sufficient e-foldings in the standard inflation era. They also analyse the cosmological perturbations generated when slow-roll is **violated after super-inflation** and show that loop quantum effects can in principle leave an **indirect signature** on the largest scales in the CMB, with some loss of power and running of the spectral index. **Shinji Tsujikawa et al 2004 Class. Quantum Grav 21 5767 doi:10.1088/0264-9381/21/24/006 Loop quantum gravity effects on inflation and the CMB Quantum geometry from quantum information** is the study of **Florian Girelli and Etera R Livine .**Loop quantum gravity defines the quantum states of space geometry as spin networks and describes their evolution in time. We reformulate spin networks in terms of harmonic oscillators and show how the holographic degrees of freedom of the theory are described as matrix models. This allows us to make a link with non-commutative geometry and to look at the issue of the semi-classical limit of loop quantum gravity from a new perspective. This work is thought of as part of a bigger project of describing quantum geometry in quantum information terms.

**Florian Girelli and Etera R Livine 2005 Class. Quantum Grav 22 3295 doi:10.1088/0264-9381/22/16/011 Reconstructing quantum geometry from quantum information: spin networks as harmonic oscillators.** Quantum geometric mechanics is the topic of study by **Dorje C. Brody Lane P. Hughston** establishing a clear relation between the quantum mechanics and quantum geometry and also the entangled states. The manifold of pure quantum states can be regarded as a complex projective space endowed with the unitary-invariant Fubini–Study metric. Note how differential gheometry concepts with Nunney's model could be entangle so that it is feasibly and viably applied to various type of model including consciousness (See Doctrine of Vibration)

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $\boldsymbol{Z}_7^2$ and $Z_7$in (A2.2):

$W_7^2 (X_7^2$-3 $M_7\ A_7\ X_7$- $M_7^2 (X_7$- $Y_7) + 2\ M_7^3\ A_7$+2 $A_7^2\ M_7^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_7$=1-cos ($Z_7 T_7$).

In the absence of a delay the system is stable for $A_7 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7 > (X_7 - Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$A_7 > [\ X_7 - Y_7/2\ M_7$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_7, A_7) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_7 \le 2$ *is* of interest First consider the extreme values

$$Q_7(0, A_7) = A_7^2\ (\mathrm{X}_7 - Y_7)$$

$$Q_7(2, A_7) = \left(A_7 + 2\,\mathrm{M}_7\right)^2 (2A_7\,\mathrm{M}_7 - (\mathrm{X}_7 - Y_7))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7 M_7^2 - 3M_7 X_7\right) + W_7\left(3M_7 A_7^2 - 2A_7 X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7 + Y_7\right)$$

$$=2\,M_7^2 W_7^2 (M_7 + \mathrm{A}_7) + M_7\,A_7^2 W_7 + \left(2A_7 M_7 - X_7 + Y_7\right)(A_7 W_7 + M_7 W_7^2$$

$$+ (2- W_7)((2\,M_7\,X_7 + M_7\,Y_7)\,W_7 + A_7\,(X_7 + Y_7)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $\mathrm{W}_7$touches the $W$ –axis only at its minimum value in which case (A2.9) shows that

further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$has no minimum in the range 0≤ $W_7$ ≤2.Defining

$Q_7(W_7, A_7) = a_7 W_7^2 + c_7$ *where* $a_7\, b_7$ *and* $c_7$ Are functions of $A_7$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7 \geq b_7 \geq 0$$

And

$$b_7^2 \geq 4a_7 c_7$$

$$Case(i): Y_7 = 0$$

*Let* $Y_7 = 0$ And try the substitution A= $X_7 = /2M_7\, in$ (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_7 X_7 /$ $2M_7$) is positive for *W<2. Since* $Q_7$ increases with $A_7$(see A2.9), increasing $A_7$will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

$Case\ (ii): Y_7 = y_7\, X_7$ ,

$Let\ Y_7 = y_7\ X_7$And try the substitution $A_7 = X_7/(2 + A_7\ y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + n_7 y_7)$ we have

$$W_7^2\ y_7\left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2\left(2 - n_7 + n_7 y_7\right)\right]$$

$$-W_7\ X_7\left[\frac{X_7(1 + n_7 y_7)}{M_7^2(2 + n_7 y_7)^2} + 2\ y_7\left(n_7 - 1\right) + \frac{X_7^2(7 + y_7)}{M_7^2(2 + n_7 y_7)} = 0.\ (A2.12)\right.$$

Since $n_7$ and $A_7$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7), the\ condition (A2.77?) can\ be\ written\ as$

$[\frac{X_7(E_7-7)}{E_7^2}]^2[\frac{X_7^2}{M_7^2} - 4\ y_7\ E_7^2\frac{X_7}{M_7^2} + 4y_7(n_7 - 2)E_7^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$2\ E_7[y_7 E_7 + (y_7\ E_7(7 + y_7)(E_7 - n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - y_7 E_7(7 + y_7)(E_7 - n_7))^{7/2}]$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0 $W_7$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7-7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2-2y_7\frac{(E_7+n_7^2y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right)-4y_7(E_7-n_7)$$

And has the solution

$$\frac{X_7}{M_7^2}>\frac{E_7}{(E_7-7)}[y_7(E_7+n_7^2y_7)+[y_7^2(E_7+n_7^2y_7)^2+4y_7(E_7-7)(E_7-n_7)]^{7/2}]$$

Given $E_7>n_7$this lower limit to $X_7/M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7/M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$in the appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7<n_7$ $(ie.,2A_7M_7^2<X_7-Y_7$)then (A2.7)and (A2.8)show that $Q_7$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2}>2E_7[y_7E_7-(y_7E_7(7+y_7)(E_7-n_7)]^{7/2}]$$

And

$$E_7>n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2 > 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.

If $A_7^2 < 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.+ $(Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$ and $X_7$ and $Y_7$ by $W_7^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7 \lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]\, e^{2T_7\lambda_7} + M_7 V_7\, e^{2T_7\lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i} Z_7$ *and* separate

(A3.1)into real and imaginary parts

$$Z_7^2 + (M_7 + V_7) Z_7 \sin(Z_7 T_7) - (U_7 M_7 + Y_7) + (M_7(U_7 + V_7) - X_7)$$

$$\mathrm{X} \cos(Z_7 T_7) - 2M_7 V_7 \cos^2(Z_7 T_7) + M_7 V_7 = 0$$

$$(M_7 + U_7) Z_7 - (M_7 + V_7) Z_7 \cos(Z_7 T_7) + (M_7(U_7 + V_7) - X)$$

$$\mathrm{X} \sin((Z_7 T_7) - M_7 V_7 \sin(Z_7 T_7) = 0$$

Using (A.3.3) to eliminate $Z_7$ and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - cos\, Z_7\ T_7$ and $A_7 = U_7 - V_7$ we obtain

$$4\,M_7^2 V_7^2 W_7^4 + 2\,M_7 V_7\ W_7^3[(M_7 - V_7)^2 + A_7(M_7 + V_7) + 2(A_7 M_7 - X_7)]$$

$$+ W_7^2[M_7 - V_7)^2(2\,A_7\ M_7 + Y_7 - X_7) + (2A_7\ M_7 - X_7)^2 - A_7(M_7 - V_7)(2A_7\ M_7 - X_7)$$

$$+ M_7 V_7\,(A_7^2 + 4(X_7 + Y_7))]$$

$$+ W_7[A_7^2 + 2X_7)(2\,A_7\ M_7 - X_7) + 2A_7 Y_7(M_7 + V_7) - M_7 A_7^3] + A_7^2(X_7\ M_7 - Y_7) = 0$$

Being a quartic in $W_7$ (A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7 > 0$$

$$and$$

$$C_7 > (X_7 - Y_7)/2\,M_7$$

*recalling that* $A_7 = U_7$ and defining $C_7 = U_7 + V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0, U_7 V_7 = A_7^2(X_7 + Y_7)$$

$$Q_7\left(2, U_7 V_7\right) = \left(C_7 + 2M_7^2\right)^2 (2C_7 M_7 + Y_7 - X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**Module numbered five**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This is the condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference

between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \quad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t - T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\left(R(t)\right)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

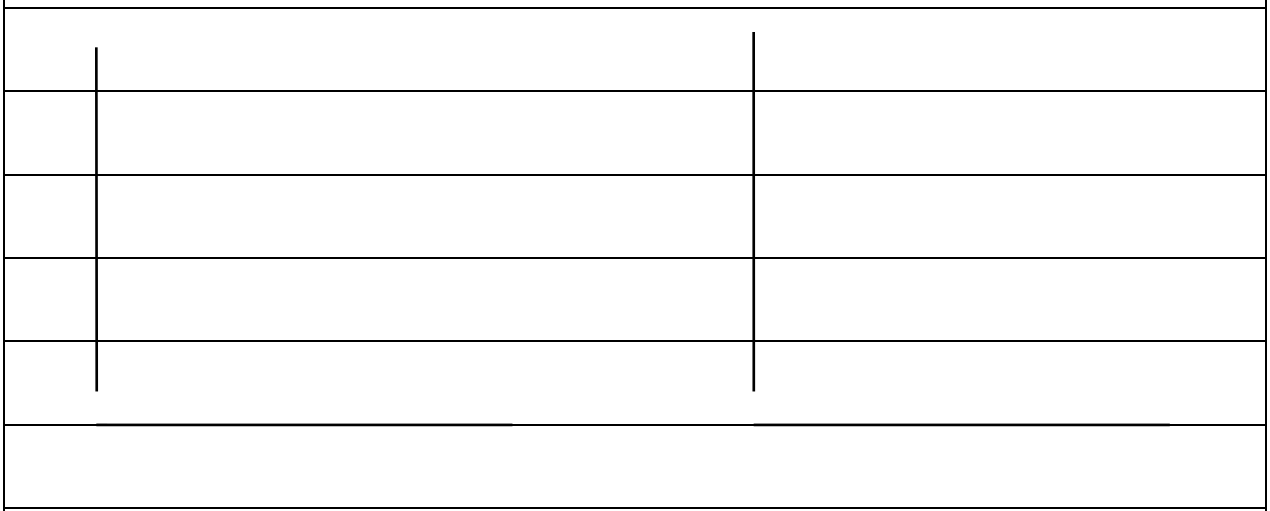
Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time

delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the



TIME DELAY

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period,

which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

As it is usually expressed outside of phenomenological texts, to undergo a conscious experience necessarily means that there is something it is like for the subject to have that experience (Nagel 1974; Searle 1992). This is obviously true of bodily sensations like pain. But it is also the case for perceptual experiences, experiences of desiring, feeling, and thinking. There is something it is like to taste chocolate, and this is different from what it is like to remember what it is like to taste chocolate, or to smell vanilla, to run, to stand still, to feel envious, nervous, depressed or happy, or to entertain an abstract belief. Yet, at the same time, as I live through these differences, there is something experiential that is, in some sense, the same, namely, their distinct first-personal character. All the experiences are characterized by a quality of mineness or for-me-ness, the fact that it is I who am having these experiences. All the experiences are given (at least tacitly) as my experiences, as experiences I am undergoing or living through. All of

this suggests that first-person experience presents me with an immediate and non-observational access to myself, and that consequently (phenomenal) consciousness consequently entails a (minimal) form of self-consciousness. To put it differently, unless a mental process is pre-reflectively self-conscious there will be nothing it is like to undergo the process, and it therefore cannot be a phenomenally conscious process. The mineness in question is not a quality like being scarlet, sour or soft. It doesn't refer to a specific experiential content, to a specific what; nor does it refer to the diachronic or synchronic sum of such content, or to some other relation that might obtain between the contents in question. Rather, it refers to the distinct givenness or the how it feels of experience. It refers to the first-personal presence or character of experience. It refers to the fact that the experiences I am living through are given differently (but not necessarily better) to me than to anybody else. It could consequently be claimed that anybody who denies the for-me-ness of experience simply fails to recognize an essential constitutive aspect of experience. Such a denial would be tantamount to a denial of the first-person perspective. It would entail the view that my own mind is either not given to me at all — I would be mind- or self-blind — or is presented to me in exactly the same way as the minds of others. There are also lines of argumentation in contemporary analytical philosophy of mind that is close to and consistent with the phenomenological conception of pre-reflective self-awareness. Alvin Goldman provides an example: [Consider] the case of thinking about x or attending to x. In the process of thinking about x there is already an implicit awareness that one is thinking about x. There is no need for reflection here, for taking a step back from thinking about x in order to examine it…When we are thinking about x, the mind is focused on x, not on our thinking of x. Nevertheless, the process of thinking about x carries with it a non-reflective self-awareness (Goldman 1970, 96). A similar view has been defended by Owen Flanagan, who not only argues that consciousness involves self-consciousness in the weak sense that there is something it is like for the subject to have the experience, but also speaks of the low-level self-consciousness involved in experiencing my experiences as mine (Flanagan 1992, 194). As Flanagan quite correctly points out, this primary type of self-consciousness should not be confused with the

much stronger notion of self-consciousness that is in play when we are thinking about our own narrative self. The latter form of reflective self-consciousness presupposes both conceptual knowledge and narrative competence. It requires maturation and socialization, and the ability to access and issue reports about the states, traits, dispositions that make one the person one is. To claim that every kind of self-consciousness is conceptual is overly cognitive. Bermúdez (1998), to mention one further philosopher in the analytic tradition, argues that there are a variety of **nonconceptual forms of self-consciousness** that are "logically and ontogenetically more primitive than the higher forms of self-consciousness that are usually the focus of philosophical debate" (1998, 274; also see Poellner 2003). This growing consensus across philosophical studies supports the phenomenological view of pre-reflective self-consciousness. (References: Stanford Encyclopedia, Post Modersnist Theories, Internet Encyclopedia of Philosophy) Korzybski pointed out that there are two ways to slice easily through life; to believe everything or to doubt everything. Both ways save us from thinking. So, from time to time you have to give yourself a Virtual Reality Check. The hallmark of postmodern philosophy has been disbelief or skepticism of all "metanarratives," or translations of reality. Postmodernism has even turned its profound skepticism on such important humanist concepts as "objective truth" and reason.Yet, for a deconstructionist postmodern society, individually we are still riddled with superstition and gullibility, and open to manipulation through our belief systems as any politician, philosopher, clergy, or salesperson will attest. Further, most people are painfully naive when it comes to even the simplest scientific understanding. Most of us don't have a clue about the fundamental nature of physical reality or our own psychological nature, and our ability to be fooled by our senses and mind, or methods of social and technological persuasion. We may think we know, but the penetration is deeper than we dare imagine. Nature rejects the naiveté that seeks absolute truth. We are beginning to realize, individually and culturally, that "realities" are all human constructions. The task becomes one of "catching ourselves in the act" of creating our own "reality" from the flow of events. Human truth is always an engagement of mind with experience.. We don't need to fear the collapse of our

personalistic belief system (the "box" we live in), nor our belief in absolute truth. **A strong desire to engage in the "quest for uncertainty" complements our anxiety that perhaps there is no absolute, objective ground to reality**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t-T_8)F_8(R_8(t-T_8) - N_8(t)M_8((R_8(t) \qquad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8} + \frac{N_8 dG_8}{dR_8} - \frac{dB_8}{dR_8} > \frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8} + \frac{dM_8}{dR_8}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_8 = \frac{N_8 G_8}{2M_8} \frac{d(8nF_8)}{dR_8} (\geq 0) \quad (8a)$$

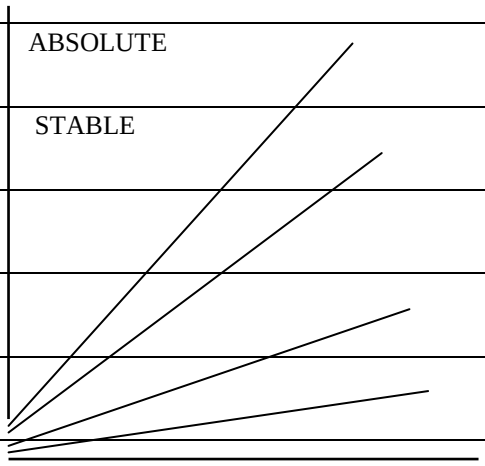
$$I_8 = -\frac{N_8 G_8}{2M_8} \frac{d(8nM_8}{dR_8} (\geq 0) \quad (8b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8 > L_8 - I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_8 > L_8 - I_8$, when $P_8^2/4 > I$ (8a)

$P_8 > L_8 - I_8 + \frac{(I_8 - P_8^2}{(L_8 - I)}$ when $P_8^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

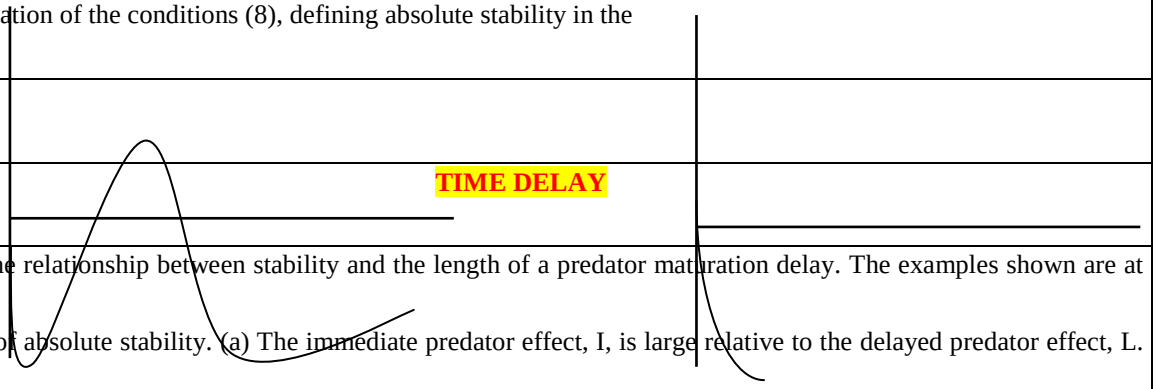



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this

rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior , we must note that the model could be applied to psychological systems too, albeit we are still not sure of the fact how the wave function collapses thereby due to perception.

Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility. He privatizes the sublime, having learned from the fate of his predecessors that the public can never be more than beautiful."(11) Notice is made here of what is indeed the telling feature. Levinas is praised, for example, for his alarm when "the social will [is] sought in an ideal of fusion . . . the subject losing himself in a collective representation, in a common ideal. . . . It is the collectivity which says 'us,' and which, turned toward the intelligible sun, toward the truth, experience, the other at his side and not face to face with him. The tellingly titled D'un ton apocalyptique contains passages that are particularly excited: These people situate themselves outside the ordinary, but they have in common this: they describe themselves as having an immediate and intuitive relation with mystery. And they want to attract, to seduce, and lead others to the mystery, through mystery. This agogic function of the leader of men, of the duce, of the Führer, of the leader, places him above the crowd that he manipulates with the aid of a small numbers of adepts joined together in a sect with a secret language, a clique or a small party with its

ritualized practices. The mystifiers pretend to have exclusive access to the privilege of a secret mystery. . . . The revelation or the unveiling of the secret is something that they jealously reserve for themselves. Jealousy is here a major characteristic. There is not narcissism and non-narcissism; there are narcissisms that are more or less comprehensive, generous, open, extended. What is called non-narcissism is in general but the economy of a much more welcoming, hospitable narcissism, one that is much more open to the experience of the other as other. I believe that without a movement of narcissistic reappropriation, the relation to the other would be absolutely destroyed; it would be destroyed in advance. The relation to the other--even if it remains asymmetrical, open, without reapporpriation--must trace a movement of reappropriation in the image of oneself for love to be possible, for example. Love is narcissistic. Beyond that, there are little narcissisms, there are big narcissisms. . . .(23) **Justice of the Pieces Deconstruction as a Social Psychology Douglas Collins** Sabine Wilke cogently argues for the rise of the aesthetic dimension in philosophical writing by showing how the metaphorical and poetic style of Heidegger, the principles of phrase configurations with Adorno, or the topographic writing of derrida replace through their very form something what traditionally would have been subject to discursive exposure The concern shifts from telling to showing, showing first and foremost that something remains untold. An exemplary book in this aesthetic drive is Avital Ronell's Telephone Book where the playful layout and explicit break with traditional design interrupt the sequential logic of the paragraphs thereby performing the very argument at stake: open up new circuits of meaning and signification. The engagement with form and publishing formats is only consequential since media theory as exposed by Walter J. Ong, Eric A. Havelock, Parry Miman or Marshall McLuhan besides many others made it clear that different communication technologies condition the mode of thought of a given culture The fractals of Benoit Mandelbrot or the butterfly shaped Lorenz attractor are the popular icons of this turn in the sciences. For communication studies and philosophy, Vilém Flusser was one of the pioneer thinkers reflecting the rise of what he called “techno-images.” The photograph is the paradigmatic case which condenses the whole crises of Western

thought and culture on its grainy surface. It subverts the linear logic of writing.6 Flusser's central diagnosis detects a shift from the grand narratives associated with the literate mind and Modernity towards creating and envisioning the world through image based practices. Writing himself in a succinct, visually inspired style, Flusser contents that the web of meaning, spun throughout the centuries by our textual practices, has become transparent. We discover the void between the letters through which the world escapes: "We are alienated from the world of text, precisely because we see through them as our own product."7 As a consequence we are confronted with a world of dissociated facts and information bits. Flusser hails this fragmented world as the advent of the zero dimensional point-universe that will provide the new raw material for our image based practices. Techno-images are nothing else than a conscious, inherently artistic drawing-together of point-elements as exemplified first and foremost by the silver-grains of the photograph but also by more recent pixel based techno-images. Similar to blueprints, designs or statistical curves, the weather map for example visualizes and contracts a tremendous amount of meteorological data from widely distributed stations into a single image of high- and low pressure

areas, temperatures, wind speeds, etc. Techno-images "compute" in this sense the universe of unimaginable, abstract information-bits into a concrete surface – the image. Echoing Paul Feyerabend, these image practices have to be understood as creative practices where the truth of scientific models has been replaced by its aesthetic quality **Aesthetics, Writing, Networked Computers Jörg Muller Dissertation Submitted to the Division of Media and Communications of The European Graduate School in Candidacy for the Degree of Doctor of Philosophy**

**Lothar Schäfer (1) and Sisir Roy (2) in** Quantum Reality, the Importance of Consciousness in the Universe, the Discovery of a Non-Empirical Realm of Physical Reality, and the Convergence with Ancient Traditions of Indian and Western Philosophy investigate metaphysical aspects that include 1) the discovery of a non-empirical part of physical reality in a realm of potentiality; 2) the emanation of the empirical world out of a realm of non-material forms; 3) the discovery that the nature of physical reality is that of an indivisible Wholeness – the One; and 4)

panpsychism: the possibility that the One is aware of its processes like a Cosmic Consciousness. The convergence of powerful traditions of seemingly disparate cultures is particularly important to point out in the present process of globalization, when a unifying view is needed to avoid controversy and conflict. To quote them in extenso : This is the nonlocality of the quantum world. Menas Kafatos and Robert Nadeau (1990) have drawn a remarkable conclusion from this phenomenon: If reality is nonlocal, the nature of the universe is that of an undivided wholeness. Because our consciousness has emerged from this wholeness and is part of it, it is possible to conclude that an element of consciousness is active in the universe: a Cosmic Consciousness. Pre-reflective self-consciousness is pre-reflective in the sense that (1) it is an awareness we have before we do any reflecting on our experience; (2) it is an implicit and first-order awareness rather than an explicit or higher-order form of self-consciousness. Indeed, an explicit reflective self-consciousness is possible only because there is a pre-reflective self-awareness that is an on-going and more primary self-consciousness. Although phenomenologists do not always agree on important questions about method, focus, or even whether there is an ego or self, they are in close to unanimous agreement about the idea that the experiential dimension always involves such an implicit pre-reflective self-awareness. In line with Edmund Husserl (1959, 189, 412), who maintains that consciousness always involves a self-appearance (Für-sich-selbst-erscheinens), and in agreement with Michel Henry (1963, 1965), who notes that experience is always self-manifesting, and with Maurice Merleau-Ponty who states that consciousness is always given to itself and that the word ‘consciousness’ has no meaning independently of this self-givenness (Merleau-Ponty 1945, 488), Jean-Paul Sartre writes that pre-reflective self-consciousness is not simply a quality added to the experience, an accessory; rather, it constitutes the very mode of being of the experience: This self-consciousness we ought to consider not as a new consciousness, but as the only mode of existence which is possible for a consciousness of something (Sartre 1943, 20 [1956, liv]). The notion of pre-reflective self-awareness is related to the idea that

experiences have a subjective 'feel' to them, a certain (phenomenal) quality of 'what it is like' or what it 'feels' like to have them.

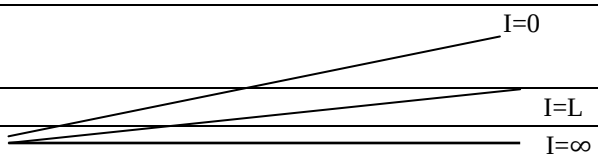
Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

TOTAL

ABSOLUTELY STABLE

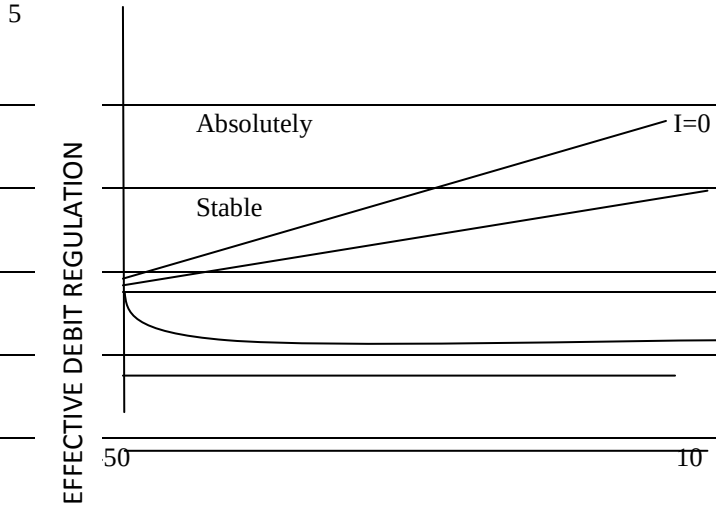



**DELAYED DEBIT EFFECT**

FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

## Utility Analysis And neural Network

Caroline Williams investigated affective density of the political and its effect on our understanding of political subjectivity. Taking up Spinoza's challenge to think about affect beyond corporeal embodiment, he argues that there is a modality of affectivity that cannot simply be inscribed within the borders of subjectivity. theorising affect as an impersonal force anchored in a relational ontology that gives due recognition to the circulation of affects, as well as to their ambivalent structure in creating sites of identification, and he utilises this ontology to reflect on the dynamic of the political and the shape of political subjectivity. Williams argues that Spinoza's philosophy (through ideas of conatus and imagination) offers the conceptual resources to reconfigure the composition of affective subjectivity as a transindividual social bond and as an unconscious dynamic of ethico-political existence. **Subjectivity (2010) 3, 245–262. doi:10.1057/sub.2010.15 Affective processes without a subject: Rethinking the relation between subjectivity and affect with Spinoza Caroline Williams** Phenomenological Time bears ample evidence and infallible observatory and apodictic evidence albeit with a sense of chagrined circumvention and defeated disgruntlment circularity of both psychological and phenomenological interpretations of art which are sometimes insightful, but most often work on the logic of cyclical time and unavoidable drives and representations played out as a ‘natural’ order: where ‘men can’t avoid being boys’ a circularity which valorizes a helplessness/truthfulness conundrum. One begins with a classical phenomenological fallacy with Duchamp’s Fountain, does the urinal keep returning or do we keep returning to the urinal? Put in this way, it is a male question, something that phenomenology conveniently conceals when positing the “I”, or indeed, the “we”. But we must be careful not to simply assume that the phenomenological first person is white, heterosexual, able-bodied and male, providing a foundation upon which knowledge should be based as a universal condition. Deleuze’s second critique of time is linked to Kant. This model of time extends not only to phenomenology but to the phenomenological interpretation

of art and is seen as a straight line. In the Critique of Pure Reason, Kant cuts up circular time and reassembles it as a series of sensory experiences which memory processes and which constitute subjectivity. Whereas the circular model of time (and interpretation) inscribes fatalism into the subject and the world, the phenomenological method is in danger of solipsism, where the subject constitutes the world. Much of embodied philosophy and enactivist theories of cognition go down this line with a plethora of art historical analyses by the likes of Rosalind Krauss, Michael Fried and others taking the lead from Merleau-Ponty. These approaches successfully challenge the **homunculus model of a Cartesian mind-body** dualism where there is a supervisor in the mind who processes experience and even consciousness and to whom all representations are directed. But phenomenology is mainly concerned with perceptual consciousness, the processing of colors, shapes, forms, motion in art and the world. The senses and sensorimotor processes are the fundamental pathways by which the interior is connected to the exterior in a chiasm of self-constitution in the world. The problem with this is that we explore the art through our sense perceptions, and what does art reveal to us?—our sense perceptions! Such solipsistic approaches have real problems relating complex and abstract conceptual production to bodily and sensorimotor contingencies, and they nearly always posit the subject as the site and object of discovery in the world. To be fair, phenomenology has moved on to rather more sophisticated models of intersubjectivity, taking into account the specificities of different kinds of bodies and ways of being beyond the white, male heterosexual as the common denominator, yet many naïve approaches to art perpetuate the notion that perceptual, rather than conceptual processes and evolutionary neural patterns are primary, fundamental ways to understand the art experience In fact, these aspects are only part of art production and reception, aspects which are often filtered out or made known to us in order for us to parody such trigger responses, as in much of Duchamp's work, which holds in contempt such retinal approaches to art. If we take the **subject of abstract expressionism** in the 1950s or later minimalist art, art history, teaching and interpretation still insists on putting the act of perceiving as fundamental to the constitution not only of the phenomenological first

person but also of art itself. The male act of usually standing up, which consists of directing a stream with the hand(s) into a container or on a flat surface, sand or snow in order to leave patterns and traces (the usual child's prank) underlies the logic of macho 1950s Abstract Expressionist culture, especially that of Jackson Pollock, so effectively parodied by Carolee Schneeman's Vagina Scroll, 1975. Here, painting or writing is a kind of a discharge which invests the **object and space with** the phenomenological "I", disciplined by existential struggle and reflected back by the traces, signatures and gestures left behind in the work of art, rediscoverable as an element in the making of other subjectivities.It is precisely this kind of solipsism which is blind to the art event as something well in excess of re-establishing habits of subjective self-identity through the production and interpretation of art. The attempt to reinforce a phenomenological identity upon the event by personalizing sense experience is unable to see art as something which remodels the phenomenological 'I' and neural plasticity itself. We learn new technologies and redraw our neural pathways through such new and spontaneous experiences which draw us into intersubjective worlds. Duchamp's Fountain and Deleuze's Repetition and Différence Grégory Minis sale (I have made some changes) The relation of different groups to deconstruction is settled by its relation to groups. An American enthusiast--one who has had a role in the development of our affection for Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^{\infty} \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$then given the constraints of $f(T_8)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8 \lambda_8 e^{-T_8 \lambda_8} + [(U_8 - V_8) M_8 + Y_8] - [((U_8 - V_8) M_8 -] e^{-T_8 \lambda_8} = 0$$

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$\boldsymbol{Z}_8^2 + M_8 Z_8 \sin(Z_8 T_8) - (A_8 M_8 + Y_8) + (A_8 M_8 + X_8)\cos(Z_8 T_8) = 0$$

$$(A_8 M_8) Z_8 - M_8 Z_8 \cos(Z_8 T_8) + (A_8 M_8 - X_8)\sin(Z_8 T_8) = 0$$

Where $A_8 = U_8 - V_8$Using (A2.3), substitute for $\boldsymbol{Z}_8^2$ and $Z_8$in (A2.2):

$W_8^2 (X_8^2 - 3 M_8 A_8 X_8 - M_8^2 (X_8 - Y_8) + 2 M_8^3 A_8 + 2 A_8^2 M_8^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_8$=1-cos $(Z_8 T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8 / 2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_8 \le 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2\ (\mathrm{X}_8 - Y_8)$$

$$Q_8(2, A_8) = \left(A_8 + 2\,\mathrm{M}_8\right)^2 (2A_8\,\mathrm{M}_8 - (\mathrm{X}_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2\left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8\left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8\left(2X_8 + Y_8\right)\right)$$

$$+2A_8\left(X_8 + Y_8\right)$$

$$=2\,M_8^2 W_8^2 (M_8 + \mathrm{A}_8) + M_8\, A_8^2 W_8 + \left(2A_8 M_8 - X_8 + Y_8\right)(A_8 W_8 + M_8 W_8^2$$

$$+ (2- W_8)((2\,M_8\,X_8 + M_8\,Y_8)\,W_8 + A_8\,(X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $\mathrm{W}_8$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in *A* will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range $0 \le W_8 \le 2$.Defining

$Q_8\left(W_8, A_8\right) = a_8 W_8^2 + c_8$ *where* $a_8\, b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \square - b_8 \square 0$$

And

$$b_8^2 \square 4a_8 c_8$$

$$Case(i): Y_8 = 0$$

$Let\ Y_8 = 0$ And try the substitution A= $X_8 = /2M_8\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8 /$ $2M_8$) is positive for *W<2. Since* $Q_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\ X_8$ ,

$Let\ Y_8 = \ y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 + \ n_8 y_8)$ we have

$$W_8^2\ y_8\left[X_8 n_8 \frac{(1 + n_8 y_8)}{(2 + n_8 y_8)} + M_8^2\left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8 X_8 \left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2}+2y_8\left(n_8-1\right)+\frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)}\right]=0.\ (A2.12)$$

Since $n_8$ and $A_8$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8=\left(2+n_8 y_8\right), the\ condition (A2.88?) can\ be\ written\ as$

$$\left[\frac{X_8(E_8-8)}{E_8^2}\right]^2\left[\frac{X_8^2}{M_8^2}-4y_8 E_8^2\frac{X_8}{M_8^2}+4y_8(n_8-2)E_8^3\right]$$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8/W_8^2)$ can be expressed as

$$2E_8\left[y_8 E_8+\left(y_8 E_8(8+y_8)(E_8-n_8)\right)^{8/2}\right]>\frac{X_8}{M_8^2}>2E_8\left[y_8 E_8-y_8 E_8(8+y_8)(E_8-n_8)\right)^{8/2}]$$

Thoughts on the Social and Political Implications of Correlationism .Lion Mirror Truth be told, as my thought has evolved the issue of correlationism had fallen off the radar for me. Somehow the debate had come to seem too “philosophical” to me, too “scholastic”, too remote from what interests me: understanding why social assemblages are organized as they are, how power functions in social assemblages, and what we might do to address that power and change things. Somehow the question of whether or not we can get out of the correlation between thinking and being just came to seem remote from these sorts of issues. Somehow it seemed too epistemological. ASIDE: Numerous discussions over the years have led me to believe that the debate over correlationism is poorly understood (or maybe I just don’t understand it). On countless occasions I’ve heard people say “of course we must relate to things in order to know them.” Well yeah, of course! I don’t think this is what the critic of correlationism is getting at. It seems to me that correlationism is something more robust than the theses that we must relate to something to know it.

Correlationism instead seems to require the theses that thought and being are indiscernible. Put more concretely, the correlationist is someone who argues that we either a) can never tell whether being is merely a construction of our thought (weak correlationism), or b) who argues that thought actually constructs being (strong correlationism). In other words, correlationism is another name for idealism. One can hold that we must relate to something in order to know it without being a correlationist. As an aside I should also add that I am a correlationist about some things. For example, I think money is something constructed by society and is therefore a strong correlationist when it comes to money. At any rate, for a long time I'd become rather indifferent to debates about correlationism and philosophies of access. I had learned the lessons of speculative realism– which I could have also learned, I think, from Deleuze and Guattari, the new materialist feminists, actor-network theorists such as Latour –and had moved on. However, occasionally you come across a tone of phrase that pitches something in a different light. In The Cut of the Real, Katerina Kolozova writes, …the political problem of contemporary philosophy identified by the 'new realists' is, in fact, the product of a more fundamental epistemic problem. In his book After Finitude, Quentin Meillassoux calls this problem 'correlationism' and identifies it as an essentially post-Kantian legacy, which continues to dominate and limit philosophy. As a matter of fact, correlationism lies at (eb) the heart of postmodern theory and consists in the premise that thought can only 'think itself,' that the real is inaccessible to knowledge and human subjectivity, and that there is nothing but discursive constructs that fully determine thinking and that are meth0dologically accounted for all the way down. (1 – 2) Thought thinking only itself: Thought only encountering itself. In the jargon of postmodern and poststructuralist lingo, this would be the thesis of infinite semiosis, where signs ("thoughts") only ever relate to other signs. Within this framework, discursivity comes to be the hegemonic framework defining all of being. At the level of politics and social theory more generally, if the correlationist thesis is true the

consequences are clear: all social phenomena are discursive and all solutions to social and political problems will be discursive. The sole sphere of the political will be the discursive and all questions of politics will be questions of speech-acts and interpretation. The problem here is not that many theorists recognize that the discursive and semiotic plays an important role in the social and the political. It does and I've repeated this tirelessly. The problem is with what happens when thought or the semiotic becomes a hegemon, an "all", foreclosing our ability to recognize other forms of power. What I've wanted to say is that not all power functions discursively. In my last post and elsewhere I spoke of some other forms of politics: Thermopolitics: The politics surrounding **energy in the form of calories and fuels such as gasoline** and coal, and how our life and our very bodies are structured by energy dependencies and by being trapped in particular distributive networks that render these forms of energy available. I'm being quite literal when I speak of energy, talking about the effects, for example, of the absence of food in certain educational environments on cognition, for example; and am generally hostile to metaphorical extensions of the concept of energy which I see as erasing the dimension of real materiality. Geopolitics: The role that features of natural and built geography such as mountain ranges, rivers, oceans, soil conditions, roads, housing design, etc., play in the form that social relations take and how they impact individual bodies. Chronopolitics: The way in which the structurations of time organize what is possible for us. For example, the structuration of the working day, how much we can say and comprehend at any given time, the impact of things like the invention of the clock, etc. **Oikopolitics: This would be the domain of political economy described so well by Marxists**. So five different types of politics: Semiopolitics (or what currently dominates critical theory), thermopolitics, geopolitics, chronopolitics, and oikopolitics. No doubt there are other sites of the political or political struggle that we could speak of, but this is a good start. Also, it should be obvious that these aren't exclusive domains, but are entangled in all sorts of important

ways. For example, something might take place at the level of semiopolitics (speech, law, rhetoric, norms, and communication) that has all sorts of impact at the level of thermopolitics. Congress might decide to cut programs that fund school meal programs. This, in turn, will have a thermodynamic impact on those students that go without the calories they need developmentally and cognitively to function in a particular way. There is an entanglement here of semiopolitical and thermopolitical domains. The young student here has been constrained both at the level of **semiotic phenomena and thermodynamic structures**. The point is that if true, semiotic intervention (speech-acts, protests, interpretations, deconstructions, etc) will not be an appropriate response to all political problems because social formations are not entirely structured by the semiotic. The child in that school does not suffer from a lack of the right signs, but from a lack of calories needed to run the engine of his thought and body. Certainly semiotic interventions might be needed to render that energy available, but it is the energy itself that is at issue and the absence of that energy that forms the spider web entangling him in his position. A correlationist perspective tends to erase this as even being a site of the political. **Larval Subjects January 28, 2014 Ontological Anarché Posted by larval subjects under Uncategorized January 27, 2014 Roden on Pluralism Over at enemy industry Roden has an excellent post up on the pluralism discussion. Check it out here. January 25, 2014 Different Senses of Pluralism and Ontology Posted by larval subjects under Uncategorized Responding to one of my comments over at his blog, Jeremy Trombley:** I have deleted restricted, reformulated some sentences .Kindly bear with me. Intention is to make the discursive enucleation consistent with the subject in question**.** Recourse to epistemology of the quantum mechanics generic and fervor of string theory is a sine qua non for the further comprehension of the theory and ontology .How would the world appear to us if its ontology was that of classical mechanics but every agent faced a restriction on how much they could come to know about the classical state? We show that in most respects

it would appear to us as quantum. The statistical theory of classical mechanics, which specifies how probability distributions over **phase space evolve** under Hamiltonian evolution and under measurements, is typically called Liouville mechanics, so the theory we explore here is Liouville mechanics with an epistemic restriction. The particular epistemic restriction we posit as our foundational postulate specifies two constraints. The first constraint is a classical analog of Heisenberg's uncertainty principle; the second-order moments of position and momentum defined by the phase-space distribution that characterizes an agent's knowledge are required to satisfy the same constraints as are satisfied by the moments of position and momentum observables for a quantum state. The second constraint is that the distribution should have maximal entropy for the given moments. Starting from this postulate, authors derive the allowed preparations, measurements, and transformations and demonstrate that they are isomorphic to those allowed in Gaussian quantum mechanics and generate the same experimental statistics. They argue that this reconstruction of Gaussian quantum mechanics constitutes additional evidence in favor of a research program wherein quantum states are interpreted as states of incomplete knowledge and that the phenomena that do not arise in Gaussian quantum mechanics provide the best clues for how one might reconstruct the full quantum theory. I: http://dx.doi.org/10.1103/PhysRevA.86.012103 **Reconstruction of Gaussian quantum mechanics from Liouville mechanics with an epistemic restriction Phys. Rev. A 86, 012103 – Published 10 July 2012 Stephen D. Bartlett, Terry Rudolph, and Robert W. Spekkens.** Deleuze in fact characterizes modes of existence, with their powers and capacities. The answer is this: Deleuze approaches modes of existence, ethically speaking, not in terms of their will, or their conscious decision making power (as in Kant), nor in terms of their interests (as in Marx, for example), but rather in terms of their drives. For Deleuze, conscious will and preconscious interest are both subsequent to our unconscious drives, and it is at the level of the drives that we have to aim our ethical analysis. Here, I would like to

focus on two sets of texts on the drives taken, not from Nietzsche and Spinoza, but rather from Nietzsche and Leibniz (Leibniz being one of the first philosophers in the history of philosophy to have developed a theory of the unconscious). The first set of texts comes from Nietzsche's great early book entitled Daybreak, published in July 1881. Nietzsche first approaches the question of the drives by giving us an everyday scenario: "Suppose we were in the market place one day," he writes, "and we noticed someone laughing at us as we went by: this event will signify this or that to us according to whether this or that drive happens at that moment to be at its height in us—and it will be a quite different event according to the kind of person we are. One person will absorb it like a drop of rain, another will shake it from him like an insect, another will try to pick a quarrel, another will examine his clothing to see if there is anything about it that might give rise to laughter, another will be led to reflect on the nature of laughter as such, another will be glad to have involuntarily augmented the amount of cheerfulness and sunshine in the world—and in each case, a drive has gratified itself, whether it be the drive to annoyance, or to combativeness or to reflection or to benevolence. This drive seized the event as its prey. Why precisely this one? Because, thirsty and hungry, it was lying in wait" (D 119). This is the source of Nietzsche's doctrine of perspectivism ("there are no facts, only interpretations"), but what is often overlooked is that, for Nietzsche, it is our drives that interpret the world, that are perspectival—and not our egos, not our conscious opinions. It is not so much that I have a different perspective on the world than you; it is rather that each of us has multiple perspectives on the world because of the multiplicity of our drives—drives that are often contradictory among themselves. "Within ourselves," Nietzsche writes, "we can be egoistic or altruistic, hard-hearted, magnanimous, just, lenient, insincere, can cause pain or give pleasure" (Parkes, pp. 291-292). We all contain such "a vast confusion of contradictory drives" (WP 259) that we are, as Nietzsche liked to say, multiplicities, and not unities. Moreover, these drives are in a constant struggle or combat with each other:

my drive to smoke and get my nicotine rush is in combat with (but also coexistent with) my drive to quit. This is where Nietzsche first developed his concept of the will to power—at the level of the drives. "Every drive is a kind of lust to rule," he writes, "each one has its perspective that it would like to compel all the other drives to accept as a norm" (WP 481). DELEUZE AND THE QUESTION OF DESIRE: TOWARD AN IMMANENT THEORY OF ETHICS Daniel W. Smith. **In on a Possible Physical Metatheory of Consciousness Miroljub Dugic et al** show that the modern quantum mechanics, and particularly the theory of decoherence, allows formulating a sort of a physical metatheory of consciousness. Particularly, the analysis of the necessary conditions for the occurrence of decoherence, along with the hypothesis that consciousness bears (more-or-less) well definable physical origin, leads to a wider physical picture naturally involving consciousness. This can be considered as a sort of a psycho-physical parallelism, but on very wide scales bearing some cosmological relevance. Open Systems & Information Dynamics Volume 9, Number 2, 153 (2002) arXiv: quant-ph/0212128. On a broader perspective and background, the theoretical computation of the universe is justified.IN physics and cosmology, digital physics is a collection of theoretical perspectives based on the premise that the universe is, at heart, describable by information, and is therefore computable. Therefore, according to this theory, the universe can be conceived of as either the output of a deterministic or probabilistic computer program, a vast, digital computation device, or mathematically isomorphic to such a device. Digital physics is grounded in one or more of the following hypotheses; listed in order of decreasing strength. The universe or reality: is essentially informational (although not every informational ontology needs to be digital); is essentially computable (the pancomputationalist position); can be described digitally; is in essence digital; is itself a computer (pancomputationalism); is the output of a simulated reality exercise. Following Jaynes and Weizsäcker, the physicist John Archibald Wheeler wrote the following: [...] it is not unreasonable to imagine that

information sits at the core of physics, just as it sits at the core of a computer. (John Archibald Wheeler 1998: 340) It from bit. Otherwise put, every 'it'—every particle, every field of force, even the space-time continuum itself—derives its function, its meaning, its very existence entirely—even if in some contexts indirectly—from the apparatus-elicited answers to yes-or-no questions, binary choices, bits. 'It from bit' symbolizes the idea that every item of the physical world has at bottom—a very deep bottom, in most instances—an immaterial source and explanation; that which we call reality arises in the last analysis from the posing of yes–no questions and the registering of equipment-evoked responses; in short, that all things physical are information-theoretic in origin and that this is a participatory universe. (John Archibald Wheeler 1990: 5) David Chalmers of the Australian National University summarised Wheeler's views as follows: Wheeler (1990) has suggested that information is fundamental to the physics of the universe. According to this 'it from bit' doctrine, the laws of physics can be cast in terms of information, postulating different states that give rise to different effects without actually saying what those states are. It is only their position in an information space that counts. If so, then information is a natural candidate to also play a role in a fundamental theory of consciousness. We are led to a conception of the world on which information is truly fundamental, and on which it has two basic aspects, corresponding to the physical and the phenomenal features of the world. (Wikipedia). **Participatory consciousness forms the bastion, pillar, post, stylobate and fulcrum of the redressal measures and corrective steps taken for circumvention of many problems, identified and unidentified, spread widest commonalty. Anne Baring's sententious and pithy article will be good to have a timely outlook.** I would like to open my contribution to this meeting with the words of Ken Wilber that I quoted in the first Millennium Symposium: 'The secret impulse of life is towards greater consciousness. Maybe the evolutionary sequence is from matter to body, to mind, to soul, to spirit - each transcending and including, each with a greater depth and greater

consciousness and wider embrace' Ken Wilber, and with these from Rupert Sheldrake's contribution to the same meeting: "It is possible to make quite a good case that the sun could be conscious and if the sun, then why not the stars? Why not the galaxies? Why should there not be galactic minds? And if they communicate with each other then the distances are far too great for light to be the medium of transmission. We are into inter-galactic telepathy." After many decades of study, I know of only one diagram in all religious traditions that can adequately explain these statements; that can give us both a blue-print of the many levels of consciousness and a route-map to help us reach beyond our present level of understanding. I would like to introduce a diagram of the four worlds of Kabbalah to illustrate the theme that spirit has brought these four worlds into being and that we, living in the fourth and manifest world know nothing of the existence of the other three invisible ones. It is the clearest diagram I know of that can explain to us the multi-levelled structure of consciousness and our place in the great chain of being. We are indeed, as Wilbur suggests, evolving from matter to body, to mind, to soul, to spirit but where have we come from? It is possible that, as participants in the emanation of spirit into manifest life, we have also come from spirit to soul, to mind, to body, to matter. Everything, from the immense galactic energies to the minute forms of matter participates in one life, is generated from one creative spirit. There is nothing outside this life, this spirit. This diagram suggests that truth is not an article of faith but the experience of participating in the undiscovered reality of what has brought us into being. We cannot know truth until we have entered into communion with that reality. Continuing in similar vein **Michel Bitbol dilates upon the dialectic deliberation and polemical argumentation between** hermeneutists and eliminativists. When he formulated the program of Neurophenomenology, Francisco Varela suggested a balanced methodological dissolution of the “hard problem” of consciousness. He shows that his dissolution is a paradigm which imposes itself onto seemingly opposite views, including materialist approaches. I also point out that

Varela's revolutionary epistemological ideas are gaining wider acceptance as a side effect of a recent controversy between hermeneutists and eliminativists. Finally, he emphasizes a structural parallel between the science of consciousness and the distinctive features of quantum mechanics. This parallel, together with the former convergences, point towards the common origin of the main puzzles of both quantum mechanics and the philosophy of mind: neglect of the constitutive blindspot of objective knowledge. Phenomenology and the Cognitive Sciences Phenomenology and the Cognitive Sciences **Phenomenology and the Cognitive Sciences 2002, Volume 1, Issue 2, pp 181-224 Science as if situation mattered Michel Bitbol**

**ABSOLUTE STABILITY**

Hence if $X_8 / M_8^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_8$ ) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8 - 8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2 - 2y_8\frac{(E_8 + n_8^2 y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right) - 4y_8(E_8 - n_8)$$

And has the solution

$$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8 - 8)}[y_8(E_8 + n_8^2 y_8) + [y_8^2 (E_8 + n_8^2 y_8)^2 + 4 y_8 (E_8 - 8)(E_8 - n_8)]^{8/2}]$$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2E_8[y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSOILIDATION**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + (M_8 + U_8)\lambda_8 - (M_8 + V_8)\lambda_8 e^{2T_8\lambda_8} + (M_8 U_8 + Y_8)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8\, e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + (M_8 + V_8) Z_8 \sin(Z_8 T_8) - (U_8 M_8 + Y_8) + (M_8(U_8 + V_8) - X_8)$$

$$\mathrm{X}\cos(Z_8 T_8) - 2M_8 V_8 \cos^2(Z_8 T_8) + M_8 V_8 = 0$$

$$(M_8 + U_8) Z_8 - (M_8 + V_8) Z_8 \cos(Z_8 T_8) + (M_8(U_8 + V_8) - X)$$

$$\mathrm{X}\sin((Z_8 T_8) - M_8 V_8 \sin(Z_8 T_8) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 =$

$1 - \cos Z_8\, T_8$and $A_8 = U_8 - V_8$we obtain

$$4M_8^2 V_8^2 W_8^4 + 2M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8(M_8 + V_8) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2[M_8 - V_8)^2(2A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8(M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8(A_8^2 + 4(X_8 + Y_8))]$$

$$+W_8\left[A_8^2+2X_8\right)\left(2A_8M_8-X_8\right)+2A_8Y_8(M_8+V_8)-M_8A_8^3]+A_8^2\left(X_8M_8-Y_8\right)=0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8>0$$

$$and$$

$$C_8>(X_8-Y_8)/2M_8$$

$recalling\ that\ A_8$= $U_8$and defining $C_8=U_8+V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8(0,U_8V_8=A_8^2(X_8+Y_8)$$

$$Q_8\left(2,U_8V_8\right)=\left(C_8+2M_8^2\right)^2(2C_8M_8+Y_8-X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0$W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$the existence of a D-position at this point is determined by letting cos $(Z_8T_8)=-1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all $T$, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**ABSOLUTE STABILITY IN PREDATOR- PREY MODELS**

**APPLIED TO ACCETUATION ATTRITION MODELS**

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability

criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The Percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending

upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\left[F\left(R(t)\right) - M\left(R(t)\right)\right]$$

(3)

$$\frac{dR(t)}{dt} = B\left(R(t-T)\right) - D\left(R(t)\right) - N(t)G\left(R(t)\right).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t-T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\left(R(t)\right)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The

regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the

**TIME DELAY**

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a dominant eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

DELAY DUE TO CONGREGATIONAL AND CONFEDERATIONAL DEBIT CREDIT TYEPES

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t-T_8)F_8(R_8(t-T_8) - N_8(t)M_8((R_8(t) \quad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8} + \frac{N_8 dG_8}{dR_8} - \frac{dB_8}{dR_8} > \frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8} + \frac{dM_8}{dR_8}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_8 = \frac{N_8 G_8}{2M_8}\frac{d(8nF_8)}{dR_8} (\geq 0) \quad (8a)$$

$$I_8 = -\frac{N_8 G_8}{2M_8}\frac{d(8nM_8}{dR_8} (\geq 0) \quad (8b)$$

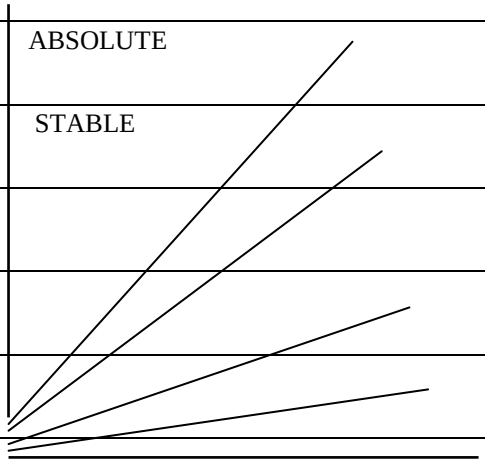
 

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8 > L_8 - I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_8 > L_8 - I_8$, when $P_8^2/4 > I$ (8a)

$P_8 > L_8 - I_8 + \frac{(I_8 - P_8^2}{(L_8 - I)}$ when $P_8^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit

(like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

See the diagram Nunney gives for predator and prey analysis. We can simulate the same for the financial system in the foregoing too.

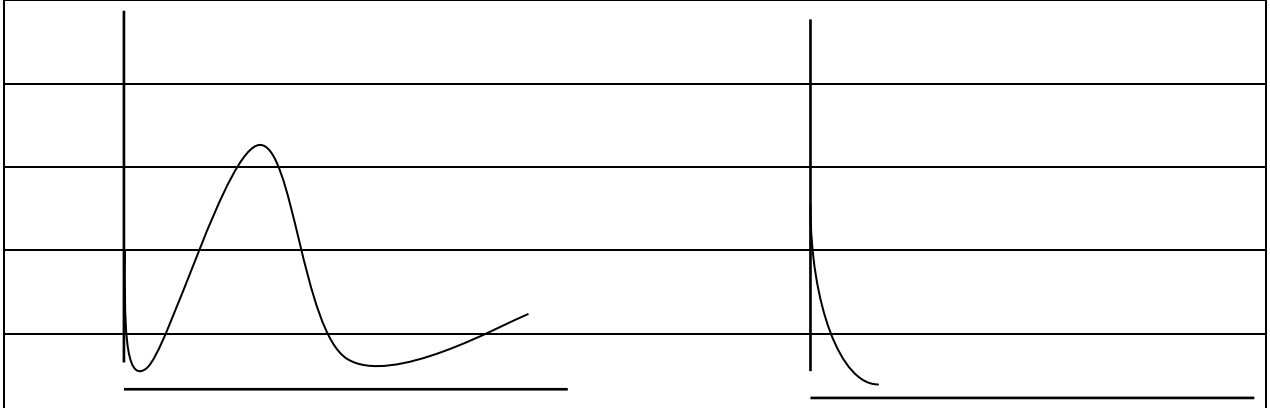
Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the

debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

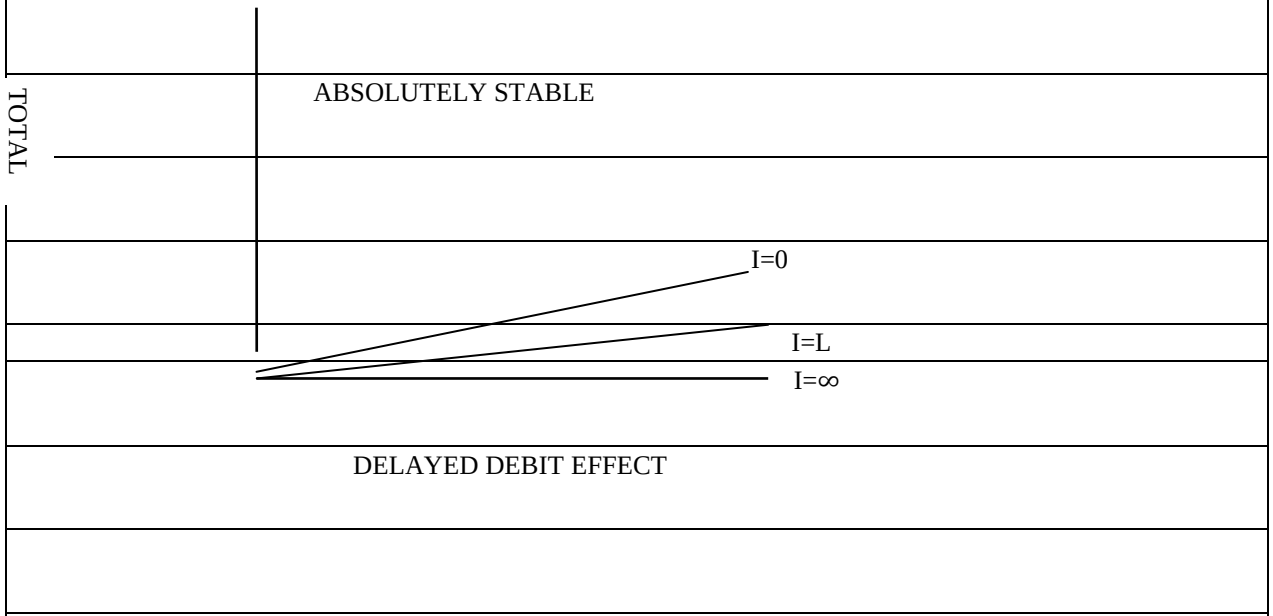
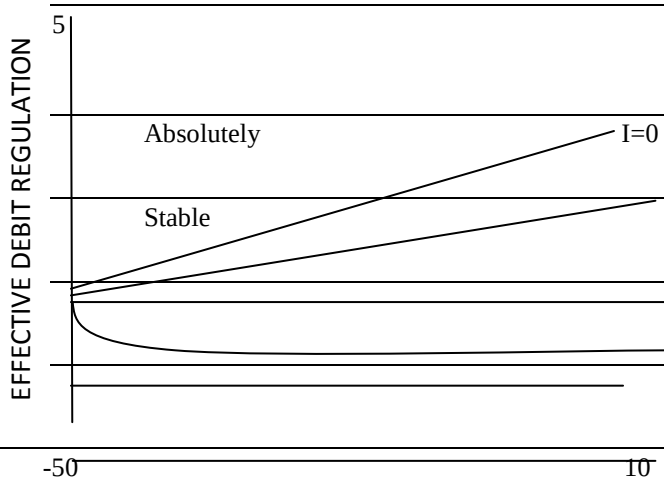



FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison?

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are

defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^\infty \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$ then given the constraints of $f(T_8)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8 \lambda_8 e^{-T_8 \lambda_8} + [(U_8 - V_8) M_8 + Y_8] - [((U_8 -$$

$$V_8 M_8 -] e^{-T_8 \lambda_8} = 0$$

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$\boldsymbol{Z}_8^2 + M_8 Z_8 \sin(Z_8 T_8) - (A_8 M_8 + Y_8) + (A_8 M_8 + X_8)\cos(Z_8 T_8) = 0$$

$$(A_8 M_8) Z_8 - M_8 Z_8 \cos(Z_8 T_8) + (A_8 M_8 - X_8)\sin(Z_8 T_8) = 0$$

Where $A_8 = U_8 - V_8$Using (A2.3), substitute for $\boldsymbol{Z}_8^2$ and $Z_8$in (A2.2):

$W_8^2 (X_8^2 - 3 M_8 A_8 X_8 - M_8^2 (X_8 - Y_8) + 2 M_8^3 A_8 + 2 A_8^2 M_8^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_8$=1-cos $(Z_8 T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8 / 2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_8 \le 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2\ (\mathrm{X}_8 - Y_8)$$

$$Q_8(2, A_8) = \left(A_8 + 2\,\mathrm{M}_8\right)^2 (2A_8\,\mathrm{M}_8 - (\mathrm{X}_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2\left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8\left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8\left(2X_8 + Y_8\right)\right)$$

$$+2A_8\left(X_8 + Y_8\right)$$

$$=2\,M_8^2 W_8^2 (M_8 + \mathrm{A}_8) + M_8\, A_8^2 W_8 + \left(2A_8 M_8 - X_8 + Y_8\right)(A_8 W_8 + M_8 W_8^2$$

$$+ (2- W_8)((2\,M_8\,X_8 + M_8\,Y_8)\,W_8 + A_8\,(X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $\mathrm{W}_8$touches the $W$ –axis only at its minimum value in which case (A2.9) shows that further increases in *A* will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range $0 \le W_8 \le 2$.Defining

$Q_8\left(W_8, A_8\right) = a_8 W_8^2 + c_8$ *where* $a_8\, b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \square b_8 \square 0$$

And

$$b_8^2 \square 4a_8c_8$$

$$Case(i): Y_8 = 0$$

$Let\ Y_8 = 0$ And try the substitution A= $X_8 = /2M_8 in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8 /$ $2M_8$) is positive for *W<2. Since* $Q_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\ X_8$ ,

$Let\ Y_8 =\ y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 +\ n_8 y_8)$ we have

$$W_8^2 y_8\left[X_8 n_8 \frac{(1 + n_8 y_8)}{(2 + n_8 y_8)} + M_8^2\left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8 X_8 \left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2}+2y_8\left(n_8-1\right)+\frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)}=0.(A2.12)\right.$$

Since $n_8$ and $A_8$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8=(2+n_8 y_8), the\ condition (A2.88?) can\ be\ written\ as$

$[\frac{X_8(E_8-8)}{E_8^2}]^2[\frac{X_8^2}{M_8^2}-4y_8 E_8^2\frac{X_8}{M_8^2}+4y_8(n_8-2)E_8^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8/W_8^2)$ can be expressed as

$2E_8[y_8E_8+(y_8 E_8(8+y_8)(E_8-n_8))^{8/2}]>\frac{X_8}{M_8^2}>2E_8[y_8E_8-y_8E_8(8+y_8)(E_8-n_8))^{8/2}]$

**ABSOLUTE STABILITY**

Hence if $X_8/M_8^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0 $W_8$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8-8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2-2y_8\frac{(E_8+n_8^2y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right)-4y_8(E_8-n_8)$$ 0

And has the solution

$$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8-8)}[y_8(E_8 + n_8^2\, y_8) + [y_8^2\,(E_8 + n_8^2\, y_8)^2 + 4\, y_8 (E_8 - 8)(E_8 - n_8)]^{8/2}]$$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2\, E_8 [y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + (M_8 + U_8)\lambda_8 - (M_8 + V_8)\lambda_8 e^{2T_8\lambda_8} + (M_8 U_8 + Y_8)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8\, e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + (M_8 + V_8) Z_8 \sin(Z_8 T_8) - (U_8 M_8 + Y_8) + (M_8(U_8 + V_8) - X_8)$$

$$\mathrm{X}\cos(Z_8 T_8) - 2M_8 V_8 \cos^2(Z_8 T_8) + M_8 V_8 = 0$$

$$(M_8 + U_8) Z_8 - (M_8 + V_8) Z_8 \cos(Z_8 T_8) + (M_8(U_8 + V_8) - X)$$

$$\mathrm{X}\sin((Z_8 T_8) - M_8 V_8 \sin(Z_8 T_8) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 = 1 - \cos Z_8 T_8$and $A_8 = U_8 - V_8$we obtain

$$4M_8^2 V_8^2 W_8^4 + 2M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8(M_8 + V_8) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2[M_8 - V_8)^2(2A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8(M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8(A_8^2 + 4(X_8 + Y_8))]$$

$$+ W_8[A_8^2 + 2X_8)(2A_8 M_8 - X_8) + 2A_8 Y_8(M_8 + V_8) - M_8 A_8^3] + A_8^2(X_8 M_8 - Y_8) = 0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8 > 0$$

$$and$$

$$C_8 > (X_8 - Y_8)/2M_8$$

$recalling\ that\ A_8 = U_8$ and defining $C_8 = U_8 + V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8(0, U_8 V_8 = A_8^2 (X_8 + Y_8)$$

$$Q_8\left(2, U_8 V_8\right) = \left(C_8 + 2M_8^2\right)^2 (2C_8 M_8 + Y_8 - X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$the existence of a D-position at this point is determined by letting cos $(Z_8 T_8) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all *T*, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**ABSOLUTE STABILITY IN PREDATOR- PREY MODELS**

**APPLIED TO ACCENTUATION ATTRITION MODELS**

We apply model to each and every module and a concatenated model was to be presented which we procrastinate due to spatial restrictions.

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - D(R) - NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

The question is political-discursive. It's what Stengers and Latour are getting at with their concepts of diplomacy and cosmopolitics. They grant, first, that all entities exist and, second, that to say that someone's cherished idol (or whatever disputed entity they hold dear) is non-existent is a **'declaration of war'** – 'this means war,' as Stengers often says. They thus shunt onto-political discourse off of the terrain of knowledge/belief in the sense of existence/non-existence. Their basic claim seems to be that 'respect for otherness,' i.e. political pluralism, can only come from granting the entities that others hold dear an ontology, even if you don't 'believe' in them. You

are thus permitted to say 'I do not follow that god, he has no hold over me' but you are not permitted to say 'your god is an inane, infantile, non-existent fantasy, grow up.' And it's not just a question of politeness (although there's that too). The point is to grant others' idols and deities an existence – one needn't agree over what that existence entails, over what capacities that entity has or what obligations it impresses upon you as someone in its partial presence but to deny it existence entirely is to 'declare war' – to deny the possibility of civil discourse, of pluralistic co-existence. I believe that it was Richard Rorty who once quipped something like claims to reconcile realism and idealism always seem to end with a triumph of idealism. We don't, in fact, get realism through such approaches, but rather just get a pervasive anti-realism. I think this is also the problem with the "non-controversial pluralism" advocated by Stengers and Latour that Phillip defends here. Such pluralism is not realism but is, in fact, a thoroughgoing social constructivism. I think this is the central problem with Latour's argument in Irreductions (these days I regret having ever defended it). In rejecting both Enlightenment critique and what he calls "reduction" he wants to say something like "The Pentacostal really is filled with (e&eb) the Holy Spirit", that for the 19th working scientist heat really is a fluid and phlogiston really is what allows things to burn, and that for the Greek lightning really is an expression of Zeus's anger. Latour tells us that we aren't to reduce or explain away the entities posited by another group's "ontology" but are to develop explanations from within that ontology. It might not sound particularly sexy– and it certainly doesn't tell us what is worth thinking –but I can't help but believe that philosophy is the critical and reflective investigation of basic concepts that guide our investigation of the world about us, how we ought to live our lives, and what form of governance might be best. Compare two figures. A scientist might ask, what causes depression? We can very well imagine a philosopher turning around and asking the scientist, what is causality? The scientist presupposes a concept of causality in her investigations. She uses this concept in her inquiry. Now she might have a sophisticated concept of causality or she might never have thought much about causality at all, using it in the sort of colloquial and unreflective way that Plato decried when, for example, people

like Euthyphro talked about piety. A whole cascade of questions arise when we raise a question like "what is causality?" We can ask whether or not causality exists at all. We can ask how we distinguish between correlation and two events that merely accompany one another from genuine causation. This, for example, was Hume's question. But perhaps most importantly we can ask whether there is only one form of causality or many forms of causality. Is there only one-to-one causation; one cause and one effect? Is there many-to-one causation; or many events conspiring to produce an effect? Is there one-to-many causation; or one event producing a variety of different effects? We can even ask whether causality necessarily moves from past to present or whether there aren't forms of causality that move from future to past! **January 23, 2014 Pluralism and Realism:** multiple-worlds Over at Struggles Forever, Jeremy Trombley has an interesting post up on "the ontological turn" in anthropology or ethnography. I've been meaning to have a discussion with him about this as I think it's an issue many of us are struggling with. For example, the core project of The Democracy of Objects– a project which I think many have missed –is to somehow reconcile some version of social constructivism with a realist ontology capable of making room for ecology (which requires realist and materialist positions as there's a fact of the matter where global warming is concerned) as well as the role played by objective agencies in social assemblages such as technologies, infrastructure, features of geography, local climates, the growth cycles of plants and animals, waste, etc. Maybe we can try to organize some cross-blog event to discuss these issues. I certainly think they're close to the heart of Jeremy, Michael of Archive Fire, Arran James, and a host of others. As an aside, I'm beginning to realize how the different sites of the political I've been outlining– semiopolitics, thermopolitics, oikopolitics (political economy), geopolitics, eropolitics (the politics of sex and desire), biopolitics, and chronopolitics (and I'm sure there are other political sites!) –are drawing me away from traditional Marxism. Assuming that classical Marxism holds that economics or the conditions and relations of production are determinative of all other sites of the political, the various sites of the political that I've been outlining would lead to the conclusion that there is not one determinative

base of the political. This would not require committing Marx to flames, but rather of recognizing the phenomenon of overdetermination, or of a variety of different entangled sites of the political. But I digress. First, I find myself wondering what the ontological turn means in ethnography. Is it 1) the investigation of the different ontologies held/proposed by different cultures? E.g., the Aztecs believed that reality was structured in this way, while the Greeks in that way, and the ancient Chinese this way, etc? Or 2) is it an investigation of how real entities– independent of cultural beliefs –influence cultural formations? Or is it a combination of both? A position that I would favor. January 20, 2014 Model is applied to structured reality

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

**The evaporation spectrum of black holes from a local quantum gravity perspective Aurélien Barrau (Submitted on 21 May 2016**)

(3) Authors revisit the hypothesis of (e) a possible line structure in the Hawking evaporation spectrum of black holes, due to (e) non-perturbative quantum gravity effects, even arbitrarily far away from (e) the Planck mass.

(4) Authors show that this naive prediction might in fact hold in (eb)the specific context of loop quantum gravity, with (e&eb) a small departure from (e) the ideal case for (e) some low-spin transitions.

They also show that the effect is (=) neither washed out by the dynamics of (e) the process, nor by (e) existence of a mass spectrum up to (e) a given width, nor by the secondary component induced by (e) the decay of neutral pions emitted during (e&eb) the time-integrated evaporation. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Cite as: arXiv:1605.06613 [gr-qc] (or arXiv:1605.06613v1 [gr-qc] for this version

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t - T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\left(R(t)\right)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

It's very easy to treat this as an abstract, academic debate: "Are you a realist or are you an anti-realist?", as if it were just a matter of what happens to be true. But it seems to me that this debate has, in reality, always been about politics. As I outlined in my last post, we have perpetually seen how appeals to the real and natural have been used in the name of oppressive power, inscribing both the exploitation of nature and the oppression of various people in the very fabric of being itself. Theistic theology and realist ontology have perpetually been used in the name of

what Deleuze called "State Philosophies" or philosophies that ontologize contingent orders of power and privilege (e.g., "the great chain of being" used to justify patriarchy, monarchy, serfdom, poverty, etc, and appeals to nature used to justify poverty and racial inequality (The Bell Curve), patriarchy (evolutionary psychology), heteronormativity, capitalism, etc). Because arrangements of power and inequality are always contingent in the sense that there's no marked difference in the capacities of peoples, power always looks for a transcendent supplement that would provide justification through ontological necessity. Antirealism– from the Greek atomists to present –became the radical and emancipatory gesture because it revealed the lie behind all of these forms of social organization or their inherit contingency or arbitrariness. Realism, by contrast, has all too often functioned as an apologetics for arbitrary power and social organizations. Here it's worth recalling what Foucault said about science: "…Even before we know to what extent something like Marxism or psychoanalysis is analogous to a scientific practice in its day-to-day operations, in its rules of construction, in the concepts it uses, we should be asking the question, asking ourselves about the aspiration to power that is inherent in the claim to being a science. The question or questions that have to be asked are: "What types of knowledge are you trying to disqualify when you say that you are a science? What speaking subject, what discursive subject, what subject of experience and knowledge are you trying to minorize when you say 'I speak this discourse, I am speaking a scientific discourse, and I am a scientist.' What theoretico-political vanguard are you trying to put on the throne in order to detach it from the massive, circulating, and discontinuous forms that knowledge can take?" (Society Must Be Defended, 10). All of these questions hold equally for claims to something being real. What is one trying to minorize when claiming something is real? What becomes privileged? What is excluded? It is these questions that have been at the heart of the realism debates, for as Spencer-Brown taught us, every distinction has a marked and unmarked space, draws attention to something to be included and pushes something into the unconscious or the domain of the invisible, hidden, or veiled. This is above all the case with evocations of the real. However, as I've tried to show antirealism

leads to its own problems. First, so long as we exclude real beings from our ontological inventory, we are unable to fully understand how power functions ("such and such a set of cultural formations have been thoroughly debunked, yet people still live as if they believed them"). Not only do we not fully understand the sources of the problems due to too much focus on the discursive and semiotic, but we deny ourselves valuable sites of political intervention at the level of infrastructure. Second, we are prevented from addressing things such as cultural racism, such as that found in Heidegger's privileging of the West and the Greeks and Germans in particular. We do a good job addressing biological and theological racism, heteronormativity, and sexism by showing how it is a cultural construction, but when faced with racism such as Heidegger's where he argues that there's something "unique" about the Greek event, it's language, and about the German language, or Badiou/Zizek's racism with respect to the "Pauline Event", we really have no response. Here someone like Jared Diamond or Fernand Braud el is needed to explain global-geographical inequalities. Third, the tools of the cultural turn really do not provide us with the means of thinking the ecological as a site of the political. For this reason, I've tried to formulate a third way– which might be called "constructivist realism" or "constructivist naturalism" –that retains the insights of the cultural turn, while also allowing a robust place for the material. I don't claim to be original in this. I think that many such as Manuel DeLanda, Deleuze and Guattari, Stacy Alaimo, Karen Barad, Jane Bennett, etc., are up to something similar. It's a vast project that requires the work of a multitude of voices, especially given the way in which the culturalism pervades humanities. **The Relational and the Non-Relational: Notes towards an Immanent and Pluralist Theory of Meaning Posted by larval subjects under uncategorized k-bigpic I have deleted or reformulated some sentences for which I beg pardon attributable to spatial constraints** What is correlation, what is causation? What is coordination? What is coincidence? These thoughts have bothered humans since long. I have been able to excavate some reliable material from Wikipedia which I would like to share. "All views are partial and contingent – that's the lesson of pluralism, and you say as much yourself when you say that we must be attentive to ways our own

knowledge might contain superstitions. But it's not discourse that determines whether one view works and another doesn't – it's a confrontation with **non-discursive (non-human)** agencies." I wonder if part of the issue here is that we understand pluralism differently. For me this doesn't sound like pluralism at all. Rather, it just sounds like our **epistemological condition**. We [hopefully] come to understand those portions of existence we question or investigate. The reason we come to investigate them is largely contingent. Finally, the accounts of these features of existence we give can be mistaken. All of this is perfectly consistent with a monism. I'm not sure why what he outlines above is the lesson of pluralism. It seems to me that every realist knows this. It seems to me that ontological pluralism is something quite different. Ontological pluralism is the thesis that there are many different worlds inhabited by many different entities. Thus, for example, you would have one world; say that of Lucretius, that's only inhabited by atoms and their combinations. You would have another world; say the world of the Mongolian shaman, that's inhabited by spirits that do all sorts of things. The ontological pluralist is saying that these spirits are, that they exist, that they're real. This is something quite different than merely saying that there are partial and contingent points of view on the world. It's this saying that these entities are rather than that some person or groups of people believe that they are that is the nub of the issue. Now, I think part of the issue here is that there's an ambiguity in the term ontology. Ontology can be one of two things. On the one hand, ontology is a group or persons set of beliefs as to what is. Here it's trivially true that there are pluralities of ontologies and the realist readily recognizes this. This is the whole reason there are debates over ontology. Mongolian shamans have their ontology, Europeans theirs, Christian fundamentalists theirs, materialists theirs, etc. When striving to understand and communicate with others it's vital to understand these ontologies because, as rhetoricians like Burke point out, our **beliefs about what is** are among the things that motivate our action. Despite having never seen bacteria I was my hands because I believe there are bacteria and viruses on door handles and whatnot and don't want to get sick. My belief about a particular thing existing is what motivates my action. And who knows, perhaps this belief is as

superstitious as the belief that the crops failed because God was displeased with my community. On the other hand, ontology is a theory **about what is.** It is making a claim that something exists. This is where the rubber hits the road. The ontological pluralist seems committed to the thesis that every group set of beliefs about what exists is sufficient for granting the existence of those entities. This is what I find objectionable in Latour. Obviously I'm not bothered by Latour's suggestion that we should take into account the role that nonhumans like speed bumps, rivers, microbes, etc., play in the form that social assemblages take. But this is not all that Latour claims. This is precisely where the philosopher might balk. Unlike the ethnographer, the philosopher is not interested in what people believe exists, but rather philosophers– at least of the realist variant –are trying to figure out what is. In other words, the realist philosopher begins with the premise that not all of these beliefs about what is are true. So for the philosopher, recognizing that Mongolian shaman's believe in the existence of shamans would only be the first step. The next step would consist in determining whether there's good reason for thinking such entities really do exist, i.e., whether there's good reason for believing these entities have mind and culture independent reality. Lest readers think that I'm just picking on the supernatural here, we can ask similar questions about strings, subatomic particles, galaxies, etc., etc., etc. I guess the question really comes down to what exactly we mean by ontological pluralism. When we talk about ontological pluralism are we defending the thesis that people have different theories of being? If so, then ontological pluralism is trivially true. If this is what is meant, then I certainly share Jeremy's view that it's valuable to understand the different world's people believe in. Certainly when I was practicing as an analyst I didn't get in ontological debates with my patients and it was necessary to understand their theory of being or their ontology to properly attend to them. Or, when we talk about ontological pluralism are we defending the thesis that all these ontological theories are true and refer to really existing entities? That's quite a different claim and is not one I would defend or endorse. Now someone might object that "in both The Democracy of Objects and Onto-Cartography you defend the thesis that there are multiple worlds." This is true. Because I hold that not everything

is related I'm led to the conclusion that there are diverse worlds. However, I also hold that however many worlds there might be, these worlds are nonetheless composed solely of material entities. Within the framework I propose I wouldn't suggest that there's one world where there are spirits and another world where there are souls and yet another where there are only material entities. My view is that there aren't spirits or souls in any of these worlds. One of the things I keep hearing in these discussions is that somehow the realist adopts a view from nowhere. I honestly don't understand this criticism. Investigation always occurs somewhere and requires all sorts of mediations involving technologies, experiments, etc. It's that labor of gathering evidence, conducting experiments, using technologies to observe the world, etc., that gradually gives us a body of data that allows us to say there's good reason to believe that such and such a thing exists and has these powers. Another charge seems to be that the realist refuses to recognize that their claims about the world are fallible. I find this charge particularly strange because it's precisely because the realist recognizes the difference between our theories of the world or what we say about the world and the world itself that fallibility is built into the core of his position. Realism doesn't mean one holds they have special access to the world, that they know all truths, or that they have the truth in hand, only that there are truths to be known and that we can be mistaken about things. Circling **Squares Philip** has a post responding to my quandaries about how to mesh realism and pluralism. He writes: Ontologically and metaphysically the idea of realist pluralism is no longer an issue. There are (appropriately) numerous variants but the basic idea that reality is itself pluralistic is well established. Captioned expatiation is about Everett worlds and string theory

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from

being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

The von Neumann formulation leads us to conclude that the lower realms of quantum reality form a virtually limitless array of potentials that are inherently incapable of realization without the observation of a knowing entity, which we identify with consciousness (Kafatos and Nadeau, 1990). We thus have a correspondence between quantum theory, the hierarchical structure implied by the Holographic Principle, and the higher orders of experience,

leading to the full expression of the Conscious Universe. David Bohm's theory of the implicate and explicate orders involves a holographic Principle that is fully consistent with the Holographic Principle discussed here (Bohm 1980; 1986; Bohm and Hiley, 1993). According to Bohm, there is an implicate order that represents the universal, holographic subtext of reality, and which unfolds in every moment to produce the explicate order that we all observe. Thomas Germinario (2004) has equated the implicate order with (eb) the unconscious process, and the explicate order with (eb) conscious process. He emphasizes the nature of dreams within (eb) the implicate order, and the importance of the dream work for maintaining a healthy mind through (e&eb) integration of the implicate order and our daily lives. Allan Combs and Mark Holland (1990) connected the implicate order or holomovement (Bohm, 1980) with (e&eb) Carl Jung is theory of synchronicity (Jung and Pauli, 1955), with (e&eb) the implicate order providing (eb) a holographic medium through which apparently disconnected individuals become connected. The principle of synchronicity, the instantaneous connection of people and events beyond (e) the senses, has been equated with the quantum-physical principle of non-locality (Combs and Holland, 1990; Germine, 1991), and has been proposed to be the fundamental mechanism of conscious process (Germine, 1991).Applying the Holographic Principle theory of mind, subjects may be connected synchronously in the manner of two individuals conversing through cellular phones, where the radio signals are transmitted through (e&eb) a distant satellite. The difference here is that the radio signals travel at the speed of light, and thus the there is some miniscule time lapse between the sending and receiving ends. The Holographic Principle implies (eb) a much more distant and instantaneous communication pathway, with (e&eb) the analogue of the satellite being the holographic boundary of the Universe, as well as a much richer communication, potentially involving a transmission of thoughts and feelings. Our ego-consciousness seems to mask (e) the universal relatedness implied by (e) the Holographic Principle, and it is perhaps only through transcendence of (e) the ego-consciousness that the higher orders of experience can become conscious. In the phenomenon of synchronicity, there seems to be a meaningful connection between (e&eb) individuals that

breaks through the barrier of ego-consciousness. Such a connection is reported by many individuals in the course of dreams, when (e) the ego-consciousness has been suspended, at times informing (e&eb) the dreamer of something that has happened in the life of a meaningfully-connected individual. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Impassioned-intense Connectionist neuroscience, in an effort to attain relevance to clinical practice, has recently given us a theory of the relationship of the unconscious to (e&eb) consciousness which is (=) dualistic (Viamontes and Beitman, 2007), an ad hoc theory of conclusions based on (e) assumptions, the most erroneous of which are that representations of (e) information are information, and that such representations elaborate themselves through (e&eb) a process of representations of representations in neuronal circuits (Bechtel and Abrahamsen,1991). As theory, the connectionist models are lacking (e) heuristic value and **are (=) inimical to the progress** of (e) basic and clinical neuroscience. This is not to denigrate or minimize the importance of neuronal connections, but these connections are rather like the wires in (eb) a radio that plays a symphony. The radio, in and of itself, is (=) incapable of playing (e&eb) the symphony. The radio doesn't do not write the symphony, nor does it broadcast the Passionate-raging symphony over the airways. Yet, if we loosen one wire, the symphony is (=) not manifested through the radio. The brain is (=) not subject to (e) the same vulnerability, in that it has (e) parallel and redundant networks of processing, and, it is in this sense that the connectivity of the brain becomes (eb) important. Connectionism, to the extent that it accepts (e) consciousness and the (e&eb) unconscious at all, attributes them to (e) separate groups of "circuits" in the brain, with (e&eb) consciousness-processing circuits having (e) a limited but detailed ability to analyze (e&eb) information, and unconscious-processing processing more information in less detail (Viamontes and Beitman, 2007). There is an element of truth in this view in that information that has importance and meaning resonates within the brain at higher levels of experience and at higher levels of holographic recursion. This resonance involves what we know, what we feel, and what we are capable and willing to consciously realize. The Universe, as William Blake noted,

can be seen in a grain of sand. Indeed, the grain of sand requires (e) the Universe to exist. The human organism requires (e) the Self to exist. The Self is not a simple amalgam of (e) sensory experiences, as suggested in the connectionist model (Blinder, 2007). This amalgam view of the Self can be damaging when applied to (e&eb) psychotherapy, as the connectionists suggest (Blinder, 2007). The Self is ours forever and ever, from eternity to eternity. As Milton suggested, the mind is (=) its own place. Our deepest purpose is for our minds to resonate with (e&eb) the supra-conscious levels of experience, much as it is (=) purpose of the radio to play the symphony. If the symphony is not to our liking, (eb) we can turn the radio off. So, perhaps, the transpersonal consciousness can turn (e&eb) our own individual and collective radios off, or perhaps turn the volume down, if (e) we are not in harmony with the **Universal Mind**. Physics has given us many enigmas. Quantum theory itself is an enigma. But some of the enigmas seem to involve (e&eb) the limitation of (e) our own abilities of discernment. The thought experiment of (e) "Schrödinger's Cat," as it is often interpreted (e.g. Gribbin, 1984), assumes (eb) that the **cat is not conscious, and is (=) incapable of observation in the context of (e) collapse of (e) the wave function.** This is due, in large part, to the false dichotomy of (e) information and (e&eb) experience. Our observation of other animals indicates, to those that are attuned to (e) the animal mind, who can "feel" the animal mind, that, at (eb) the very least, all mammals are (=) conscious. It would seem that the great neuropsychiatrist, Stanley Cobb (1948), may have been correct in his attribution of (e) levels of consciousness to a wide variety of organisms: "...lower animals with (e&eb) no cerebrum appear to be conscious...even plants such as the sunflower that turns towards strong light may have (e) a vague awareness or warmth and comfort. There are many degrees of consciousness and it is my contention that it is integrated at (e&eb) many levels like other important functions of (e) the central nervous system."With respect to the theory of mind, recent mainstream thinking remains (eb) classical and mechanistic. Such views of mind are advocated by our most in prominent neuroscientists (e.g. Changeux, 1985; Gazzaniga, 1985) as neuroscience research attempts to explain all mental function on the basis of (e) **pure mechanism and the localized function of**

**specific areas of the brain.** Francis Crick, the distinguished Nobel Laureate, had a very influential second career as a neuroscientist. He concluded that we are basically soulless creatures, that all mental processes could be reduced to (e&eb) identifiable, neural correlates in (eb) the brain, and that, in fact, we are (=) nothing but a collection of neurons (Crick, 1995). Crick even went so far as to say that what we call the soul is (=) a group of neurons located in the prefrontal lobes of the brain. Another Nobel Laureate, Gerald Edelman (1987), has written a critically acclaimed book on "neural Darwinism," which purports that the brain is "circuitry" is constructed through a developmental process of "survival of the fittest" neurons and neuronal connections. At the same time the total purposelessness and lack of direction in evolution is expounded, again to critical acclaim, by Richard Dawkins (1986). The idea that there essentially no self or soul imbued with agency and self determination was expounded by Gilbert Ryle (1949/2002) with his parody of the "ghost in the machine," and this work continues to be influential among many scholars of the mind. Philosophers such as Patricia Churchland (1986), who coined the term "neurophilosophy," take an "eliminative materialist" view of (e) belief, free will, and consciousness. Other philosophers, such as Daniel Dennett (1991) would make (eb) experience "epiphenomenal," a by-product of (e) brain processes with (e&eb) no effect whatsoever on (eb) the function of mind. However, unlike neuroscience, which is still in its infancy, the philosophy of mind has had a long history, including such notable figures as Plato, Aristotle, Descartes, Kant, Locke, Hume, Kierkegaard, and James, to name just a few. Perhaps, as noted by Ervin Laszlo (1974): "in today's world, most of the traditional functions of (e) cognitive synthesis have atrophied (e) and are ignored and neglected." Hopefully, what we have presented here will be a step in the right direction for what Laszlo calls a "conceptual synthesis" to help fill the need for (e) meaningful engagement is such a world. The fundamental elements of such a synthesis have been described by (e) Laszlo as follows: Conceptual synthesis performs (e&eb) at least five basic functions in the guidance of affairs. They are the mystical, the cosmological, the sociological, the pedagogical or psychological, and the editorial functions. The mystical function inspires (eb)

in man a sense of mystery and profound meaning related to (e&eb) the universe and of himself in it. The cosmological function forms (eb) images of the universe in accord with (e&eb) local knowledge and experience, enabling (eb) men to describe and identify (e&eb) the structure of (e) the universe and the forces of nature. The sociological function validates supports and enforces (eb) social order, representing (eb) it in accord with (e&eb) the nature of the universe, or as the natural or right form of (e) social organization. The pedagogical or psychological function guides (eb) individuals through stages of life, teaching ways of understanding themselves and others and presenting (eb) desirable responses to life is challenges and trials. Finally, the editorial function of (e) conceptual is to define (eb)some aspects of reality as important and credible and hence to be attended to, and other aspects unworthy of serious attention. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Kant sums up this inversion and its spirit early in the Critique of Pure Reason: Up to now it has been assumed that all our cognition must conform to the objects; but all attempts to find something about them a priori through concepts that would extend our cognition have, on this presupposition, come to nothing. Hence let us once try whether we do not get farther with the problems of metaphysics by assuming that the objects must conform to our cognition, which would agree better with the requested possibility of an a priori cognition of them, which is to establish something about objects before they are given to us. This would be just like the first thoughts of Copernicus, who, when he did not make good progress in the explanation of the celestial motions if he assumed that the entire celestial host revolves around the observer, tried to see if he might not have greater success if he made the observer revolve and left the stars at rest. (B xvi) Onticology, like all variations of object-oriented ontology, is realist in its orientation. In defending a realist ontology onticology holds that the vast majorities of objects, actants, beings, or entities are independent of humans and are what they are regardless of whether any humans regard them or register them. In short, onticology rejects any anthropomorphic, idealist, or anti-realist thesis to the effect that to be is to be the correlate of mind, spirit, the body, the human, and language or otherwise. While it is certainly the case that knowledge is necessarily dependent on the object to which it relates, the converse does not hold true. Objects are not dependent on being known, regarded, perceived, or spoken

about. As such, and to put it in Aristotlean terms, knowledge is an accident of objects, not objects an accident of knowledge. As Althusser so nicely puts it, "[n]o doubt there is a relation between thought-about-the-real and thisreal, but it is a relation of knowledge, a relation of adequacy or inadequacy of knowledge, not a real relation, meaning by this a relation inscribed in that real of which the thought is the (adequate or inadequate) knowledge" (Reading Capital, 96). Althusser goes on to remark that "[t]he distinction between a relation of knowledge and a relation of the real is a fundamental one: if we did not respect it we should fall irreversibly into either speculative or empiricist idealism" (ibid.). Onticology categorically endorses Althusser's verdict. It is a fundamental necessity to distinguish between those relations that belong to the object and those that belong to knowledge. Contemporary philosophy continuously confuses these two very different sorts of relations. Naturally the question arises of how it is possible to surmount our relation to the object so as to determine whether objects themselves possess the properties we encounter in relating to objects. In other words, given that we can only ever relate to the object in relating to the object how is it possible to surmount this relation to get at the being of the object itself? Much more will have to be said about this later– and the answers will be surprising with respect to standard prejudices about realism –however, for the moment it can be said that onticology takes its epistemological inspiration from the transcendental realism of Roy Bhaskar. Among other things, Bhaskar sought to provide a transcendental grounding for the sciences. Insofar as onticology defends the thesis that the field of being is much more vast than the field of objects investigated by the natural sciences, it parts way with the thesis that the domain of being is exhausted by the domain of natural objects. However, the general form of Bhaskar's argument holds for our realist purposes. A transcendental argument seeks to elucidate the conditions under which certain acknowledged practices and forms of cognition are possible. Kant, for example, asked what must be the case for mathematical judgments to be possible. How is it both that we are able to extend our knowledge, as if by magic, through mathematical judgments and, more significantly, that these judgments are able to provide genuine knowledge of the world despite the fact that these

forms of reasoning are not based on experience? Part of Kant's argument consisted in claiming that mind imposes the forms of space and time on the data of experience. In other words, space and time are not attributes of being itself but rather of the mind that regards being. Insofar, Kant argues, as mathematics is ultimately a rumination on the nature of space and time taken in their most abstract form and insofar as the mind imposes space and time on the manifold of sensation, it thus follows that a priori judgments about the nature of spatio-temporal relations are possible that anticipate the structure of actual-space times without directly experiencing these space-times. Why? Because any manifold of sensation must necessarily be structured by these forms imposed by intuition. **Onticology–A Manifesto for Object-Oriented Ontology Part I Posted by larval subjects under Object-Oriented Philosophy My Question Is This: You Say, "Granted There Are Deeper And Deeper And More Primitive Forms Of "Knowledge" Or Abstraction Into Mechanism, But I Wouldn't Say That The Atomic Level Brings A New Epistemic World Into View." But Doesn't Any Kind Of Knowledge, However Primitive, Imply Bringing Some Kind Of Epistemic World Into View Or Proto-View?** Joel: David, Yes Exactly. My Question Is, Just How Far Down The Complexity Gradient Can The Meaning Of 'Knowledge,' Or 'Semiotics, Or Specifically Representation (Maya) Be Stretched Without Breaking? I Can Comfortably Stretch It To The Genetic Level, As The Proto-Epistemic, Where Genetic Codes 'Represent' Phenotypic Structures Or Modes, And The Whole Of Evolution Learns From Its Mistakes And Successes As That Knowledge And History Is Encoded Into And As The Recapitulation Of The Embryogenesis Of The Organism. But An Atom, In My Model, While Certainly Reacting To Its Environment From Its Own Attractor, Essence, And 'Interiority' (Prehension) Isn't Making Representations Of It To Do So. But The Key To The Power And 'Essence' Of Representation (And Simultaneously Its Limitation) Is The Capacity Of Abstraction And Choice. Evolution Does This At A Rudimentary Level (And Anthropomorphizing A Bit) With Its Sending Into The World Its 'Random' Variations, Organismic Ideas, In A Sense. Its 'Mind' Is Literally On The Outside As Our Living World And Biosphere. A 'Choice' Is Made In The Biospheric Interiority,

And Knowledge For The Evolutionary/Embryogenetic Trajectory Is Gained By The "Differential Reproductive Success" Of These Individuals Lives Through Time, As The Increasing Intelligence, Representation, And Capacity For Choice Is Directly Injected Into The Flow Through Sexual Selection. But Where Is The Encoded/Abstracted (Enfolded) Representation Of Options And A Choice Among Them Made At The Atomic Level? I Am Working With The Most Cutting-Edge And Coherent Models For The Atom (Including My Own) That I Can Find Among The Heterodoxy, And I See No Evidence Of Representation And Choice At This Level, Or Rather No Real Way To Further Stretch The Meaning Of The Terms To This Level. Sure There Is Uncertainty From Infinite Difference And Immanent Causation (See Bohm On Infinite Causation), But Without Representation There Really Is No Choice (Although Crudely We Could Call A 'Bifurcation' A 'Choice' By The System As A Whole). So I Am Simply Saying That Representation At The Atomic Level Is Entirely Enfolded And These Are Really Just Ontic Level Phenomena. The Atomic Level, In My View, Is Pre- Not Yet Proto-Epistemic, Which I Place at the Genetic Level Where We Find the First Codes And (e&eb) Primitive Representation. This Just Means That There Is A Real Distinction Between Ontic And Epistemic (And A Gradient Between Them) And We Can't Collapse The Two Into Just The Ontic-Epistemic. And Yet, With The Distinction Intact, And The Gradient To Explore, They Remain Nondual, Univocal, ONE Or Ontic. What I Would Say Exists At All Levels (My Own Self-Aware Myth About The Given) Is What I Call The "Symbiogenesis Of Subject And Object". This Is The Spinozan Essence Of Dynamic Stability And Growth, And Leibnizian Prehension. Also Called The "Nucleation Of Observability" In Spinbitz. It Is From This Nucleus Of The Observer (Fuller) In A "Point-Free Geometry" (Whitehead) Of Pointless "Points Of View" As Conditions Of Boundary That Wilber's "All Is Perspective" Finds Its Ontic Grounding. And This Symbiogenesis Is Also A Rudimentary Or Deep-Level Native Or Primitive Intelligence, Evolution At The Involution Of The Subject-Object Interface. It Is An Exploration Of Creative Learning And Direct Immanent Awareness, Unmediated (Enfolded) By Representational Forms (Unfolded). It Just Is, In Waves (Cosmic Vertebrae)

Before It Abstracts, Represents And Incarnates The Flesh Of What Could Be. Matter Is Enfolded Seed Of The Flower Of Its Abstraction. Maya Is This Depth And Recursion Of Brahma Into And Through The Boundary Conditions Of Abstract Relation. And Brahma Is The Span Of The Emergent Rigor Of The Cosmic Ergodic Spine That Opens The Recursions Into The Field Of Infinite Difference. In Immanent-Transcendent Waves, Literally And Empirically A Limit-Cycle In Ergodicity, It Enfolds And Unfolds From Anatom To Anatomy, From Infinite Intelligence Or Omnirelational Awareness Of Self, Through The Representational Depths Of Self As Self-Consciousness. Omni-Evolution, Trans-Dynamics, And The Anatom Posted On February 21, 2013A Discussion With Tom Huston, Part II The Following Is Part 2 Of An Excerpt From A Discussion Initiated By Tom Huston On His Facebook Page, Edited And Embellished For Clarity. Find Part One, Here. In This Part We Dig Into The Nature Of The "Anatom" And The "Ergodic Spine" Of An Electro-Fractal Cosmos, And Into The Trans-Dynamic Integration Between Being And Becoming, And The X-Interface−The Crossroads Of The Ontic-Epistemic (Brahma And Maya) And The Subject-Object Polarities. Univocal_Dynamics_V1 Tom Huston Is A Founding Member Of Integral Institute And A Former Editor Of Enlighten Next Magazine. His Webpage Is Tomhuston.Com. David Marshall Is A Writer And Editor Living In Chicago. No meaning, value, or normativity to be found in nature, that there is nothing **natural beings ought to be,** but that, rather, these judgments arise from us (See last post of larval subjects) which has led some to raise valuable questions about the coherence of these claims. The problem arises when the following three propositions are taken together: 1. there is nothing outside of nature.2. Beings have no intrinsic meaning, purpose, or value (in and of themselves, there's nothing they ought to be). 3. Value judgments about what beings and being ought to be arise from us and beings like us (bonobo apes, dolphins, institutions, birds of paradise, etc). The problem arises between proposition 1 and 3. How can it both be true that there is nothing outside of being and those normative judgments belong to us and other beings capable of making normative judgments, not nature? The problem arises from restricting these judgments to humans and beings capable of

making these judgments. In making such a claim it seems as if we're saying that there's something outside of nature, something that is beyond nature, thereby violating the first thesis and potentially reintroducing the nature/culture distinction. By contrast, if we say that normative judgments are the special domain of those living beings with the proper degree of sentience to make such judgments, then we seem to reintroduce the nature/culture distinction and fall back into the sorts of problems that I outlined in my last post (and that are so nicely critiqued by thinkers such as Latour). Is there a way out of this? I don't know. Great_Chain_of_Being_2_ (lighter) ASIDE: It seems to me that this is really what the debate between realism and anti-realism, realism and socio-linguistic constructivism surrounding the new materialisms and speculative realists has really been all about.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \qquad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9}+\frac{N_9 dG_9}{dR_9}-\frac{dB_9}{dR_9}>\frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9}+\frac{dM_9}{dR_9}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usance bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_9=\frac{N_9 G_9}{2M_9}\frac{d(9nF_9)}{dR_9}(\geq 0) \qquad (9a)$$

$$I_9=-\frac{N_9 G_9}{2M_9}\frac{d(9nM_9}{dR_9}(\geq 0) \qquad (9b)$$

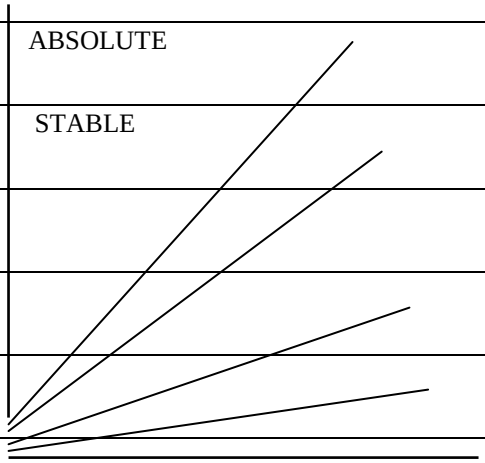
It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9>L_9-I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_9>L_9-I_9$, when $P_9^2/4>I$ (8a)

$P_9>L_9-I_9+\frac{(I_9-P_9^2}{(L_9-I)}$ when $P_9^2/4<I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Reffer couffesy Nunney

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability

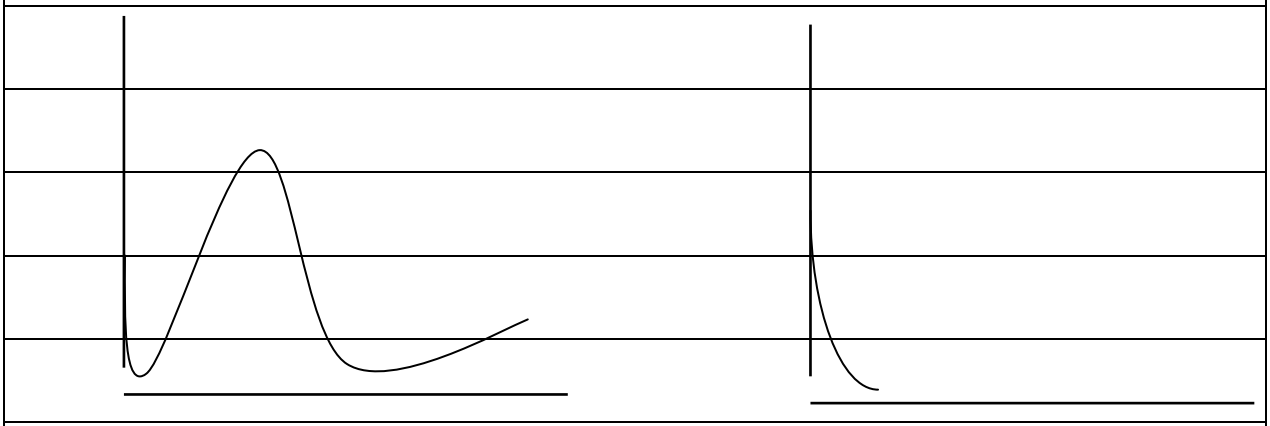
 

is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

Indeed, Nunny's model can be applied to some of the consciousness and philosophical problems. Predator prey analysis has ripped open the veneer of the fundamental concept of universe namely Debit-Credit world.

The noise of ordinary consciousness and (e&eb) beliefs drowns out (e) the signal. Unconsciousness is the background of (e) our ordinary awareness. Our organism is very much at (eb) the center of such effects. The organismic source is (=) our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is a meme-Scape. Pre-conceived concepts vie with structures, concepts with images. Like scientists who ignore assumed truths, we leapfrog over our beliefs and personality deficits, claiming idiosyncratic imagination is literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is both. We inherit the structures that make our experience what it is.

But the structure itself is “empty,” and each human culture “fills” it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us

feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through 're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns, nomenclature, and internal coherence of meaning and representation. In a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist

simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** The lower domain of experience is what we do not directly observe, the quantum realm, while the higher level is what we ordinarily observe, the classical realm. Based purely on experiment, the formulation of quantum theory initially placed the great divide between the quantum and classical domains between the measuring apparatus and the particle (Herbert, 1985). However, in the early 1930s, John von Neumann, in his rigorous mathematical treatise on quantum mechanics, found no support for such a division, compelling him to conclude that the wave for was collapsed by consciousness. Von Neumann's treatise has been called the "quantum bible" and "the most influential book on quantum theory ever written" (Herbert, 1985).

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t-T_9)F_9\big(R_9(t-T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t-T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ and the effective prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting $dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

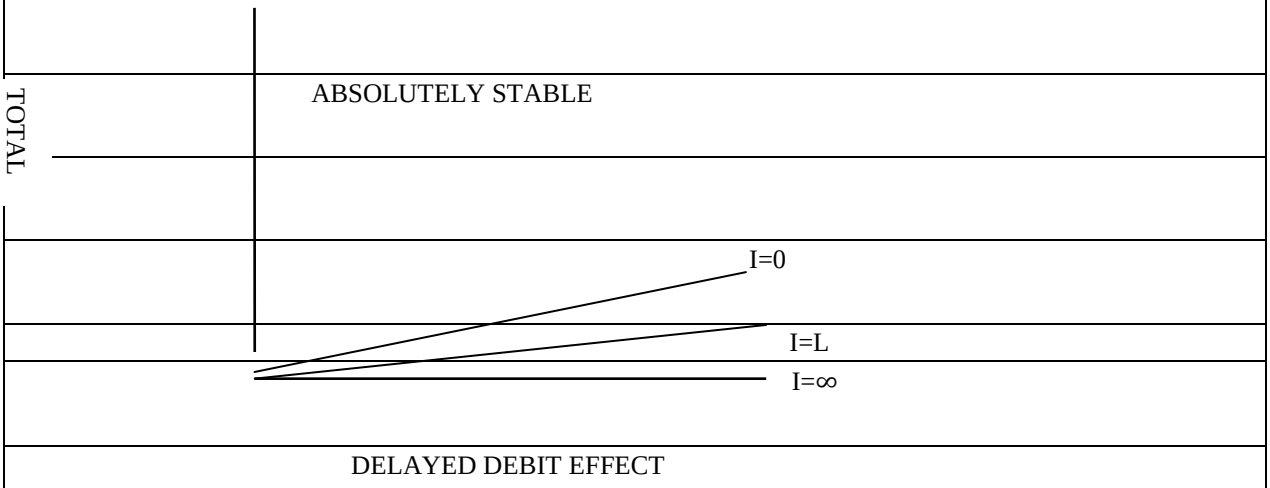



FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.

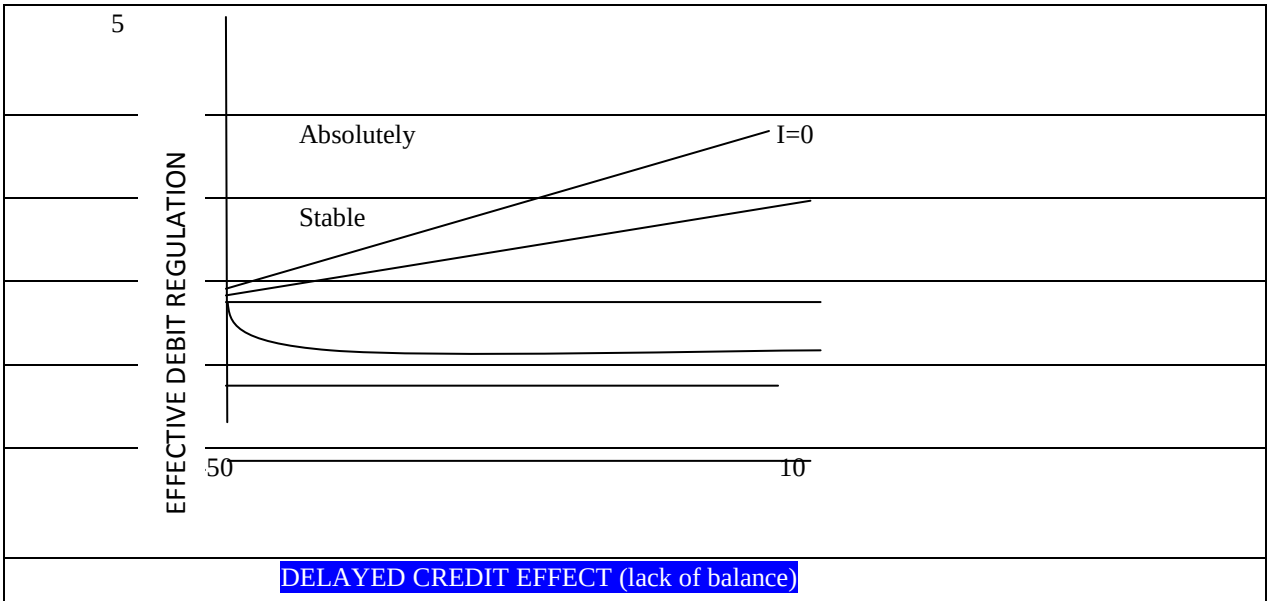



Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the

constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^\infty \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9\left(U_9 - V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9\right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$ then given the constraints of $f(T_9)$, $(A1.5)$ *can never be satisfied* Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDERATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 -$$

$$V_9 M_9 -]e - T_9\lambda_9 = 0$$

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9 Z_9 \sin(Z_9 T_9) - (A_9 M_9 + Y_9) + (A_9 M_9 + X_9)\cos(Z_9 T_9) = 0$$

$$(A_9 M_9) Z_9 - M_9 Z_9 \cos(Z_9 T_9) + (A_9 M_9 - X_9)\sin(Z_9 T_9) = 0$$

Where $A_9 = U_9 - V_9$ Using (A2.3), substitute for $Z_9^2$ and $Z_9$ in (A2.2):

$W_9^2 (X_9^2 - 3 M_9 A_9 X_9 - M_9^2 (X_9 - Y_9) + 2 M_9^3 A_9 + 2 A_9^2 M_9^2]$

$$+W\left[MA^{3}-A^{2}\mathrm{X}+2\mathrm{MA}(2\mathrm{X}+\mathrm{Y})-2X^{2}\right]+A^{2}(\mathrm{X}+\mathrm{Y})=0$$

*where* $W_9$=1-cos ($\mathrm{Z}_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (\mathrm{X}_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ \mathrm{X}_9 - Y_9/2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range 0⍰$W_9$ ⍰2 *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (\mathrm{X}_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\ \mathrm{M}_9)^2\ (2A_9\mathrm{M}_9 - (\mathrm{X}_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9M_9^2 - 3M_9X_9\right) + W_9\left(3M_9A_9^2 - 2A_9X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9 + Y_9\right)$$

$=2M_9^2W_9^2(M_9 + A_9) + M_9A_9^2W_9 + \left(2A_9M_9 - X_9 + Y_9\right)(A_9W_9 + M_9W_9^2$

$+ (2- W_9)((2M_9X_9 + M_9Y_9)W_9 + A_9(X_9 + Y_9))$,

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range 0≤ $W_9$ ≤2.Defining

$Q_9(W_9, A_9) = a_9W_9^2 + c_9$ *where* $a_9$ $b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

**Harald Atmanspacher, Hans Primas** show that such an epistemic perspective can be consistently complemented by Einstein's ontically oriented position. **Time, Quantum and Information 2003, pp 301-321 Epistemic and Ontic Quantum Realities Harald Atmanspacher, Hans Primas Robert W. Spekkens** present a toy theory that is based on a simple principle: the number of questions about the physical state of a system that are answered must always be equal to the number that are unanswered in a state of maximal knowledge. Many quantum phenomena are found to have analogues within this toy theory. These include the noncommutativity of measurements, interference, the multiplicity of convex decompositions of a mixed state, the impossibility of discriminating nonorthogonal states, the impossibility of a universal state

inverter, the distinction between bipartite and tripartite entanglement, the monogamy of pure entanglement, no cloning, no broadcasting, remote steering, teleportation, entanglement swapping, dense coding, mutually unbiased bases, and many others. The diversity and quality of these analogies is taken as evidence for the view that quantum states are states of **incomplete knowledge** rather than states of reality. A consideration of the phenomena that the toy theory **fails to reproduce, notably, violations of Bell inequalities** and the existence of a Kochen-Specker theorem, provides clues for how to proceed with this research program.DOI: http://dx.doi.org/10.1103/PhysRevA.75.032110 **Evidence for the epistemic view of quantum states A toy theory Phys Rev A 75, 032110 – Published 19 March 2007 Robert W. Spekkens In such specialized and autonomous areas, it is to the eminent and erudite authors to interpret the contradistinctions and contradictions' in the dove tailing explaining. This I feel would be yeoman service in the creation of Nature's general ledger. M. S. Leifer and O. J. E. Maroney** examine the relationship between quantum contextuality (in both the standard Kochen-Specker sense and in the generalized sense proposed by Spekkens) and models of quantum theory in which the quantum state is maximally epistemic. We find that preparation noncontextual models must be maximally epistemic, and these in turn must be Kochen-Specker noncontextual. This implies that the Kochen-Specker theorem is sufficient to establish both the impossibility of maximally epistemic models and the impossibility of preparation noncontextual models. The implication from preparation noncontextual to maximally epistemic then also yields a proof of Bell's theorem from an Einstein-Podolsky-Rosen-like argument. DOI: http://dx.doi.org/10.1103/PhysRevLett.110.120401 Received 25 August 2012Published 20 March 2013 © 2013 American Physical Society **Maximally Epistemic Interpretations of the Quantum State and Contextuality Phys. Rev. Lett 110, 120401 – Published 20 March 2013 M. S. Leifer and O. J. E. Maroney.** Epistemological, onitic and ontological aspects particularly of quantum mechanics have invited

profound thinkers to contribute their mite to the growing literature on the subject, which includes the contributions from the subterranean realm and ceratoid dualism of philosophy like that Chaos theory of Deleuze, Derrida , Kierkegaard and others. **Andrei Khrennikov** shows that the Dirac-von Neumann formalism for quantum mechanics can be obtained as an approximation of classical statistical field theory. This approximation is based on the Taylor expansion (up to terms of the second order) of classical physical variables – maps f: Ω → R, where Ω is the infinite-dimensional Hilbert space. The space of classical statistical states consists of Gaussian measures ρ on Ω having zero mean value and dispersion σ2 (ρ) ≈ h. This viewpoint to the conventional quantum formalism gives the possibility to create generalized quantum formalisms based on expansions of classical physical variables in the Taylor series up to terms of nth order and considering statistical states ρ having dispersion σ2(ρ) = h n (for n = 2 we obtain the conventional quantum formalism). **December 2005, Volume 18, Issue 7, pp 637-650 Date: 28 Nov 2005 Generalizations of Quantum Mechanics Induced by Classical Statistical Field Theory Andrei Khrennikov** It is widely accepted that consciousness or, in other words, mental activity is in some way correlated to the behavior of the brain or, in other words, material brain activity. Since quantum theory is the most fundamental theory of matter that is currently available, it is a legitimate question to ask whether quantum theory can help us to understand consciousness. Several approaches answering this question affirmatively, proposed in recent decades, will be surveyed. It will be pointed out that they make different epistemological assumptions, refer to different neurophysiological levels of description, and adopt quantum theory in different ways. For each of the approaches discussed, these imply both problematic and promising features which will be indicated. **Discrete Dynamics in Nature and Society Volume 2004 (2004), Issue 1, Pages 51-73 http://dx.doi.org/10.1155/S102602260440106X Quantum theory and consciousness: an overview with selected examples Harald Atmanspacher Copyright © 2004 Hindawi Publishing**

**Corporation. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.** Another thought provoking and investigatory article is by Peter G. Lewis et al. Perhaps the quantum state represents information about reality, and not reality directly. Wave function collapse is then possibly no more mysterious than a Bayesian update of a probability distribution given new data. **Peter G. Lewis et al** consider models for quantum systems with measurement outcomes determined by an underlying physical state of the system but where several quantum states are consistent with a single underlying state—i.e., probability distributions for distinct quantum states overlap. Significantly, they demonstrate by example that additional assumptions are always necessary to rule out such a model. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.150404 **Distinct Quantum States Can Be Compatible with a Single State of Reality Phys. Rev. Lett. 109, 150404 – Published 9 October 2012 Peter G. Lewis et al** Communication complexity of a quantum channel is the minimal amount of classical communication required for classically simulating a process of state preparation, transmission through the channel and subsequent measurement. It establishes and perpetuates a limit on the power of quantum communication in terms of classical resources. **Alberto Montina** shows that classical simulations employing a finite amount of communication can be derived from a special class of hidden variable theories where quantum states represent statistical knowledge about the classical state and not an element of reality. This special class has attracted strong interest very recently. The communication cost of each derived simulation is given by the mutual information between the quantum state and the classical state of the parent hidden variable theory. Finally, find that the communication complexity for single qubits is smaller than 1.28 bits. The previous known upper bound was 1.85 bits. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.110501 **Epistemic View of Quantum States and**

**Communication Complexity of Quantum Channels Phys. Rev. Lett 109, 110501 – Published 12 September 2012 Alberto Montina for models see elsewhere .** Stability analysis of what? Quantum state representing reality or our knowledge of reality? Does the quantum state represent reality or our knowledge of reality? In making this distinction precise, **Nicholas Harrigan, Robert W. Spekkens** are led to a novel classification of hidden variable models of quantum theory. They show that representatives of each class can be found among existing constructions for two-dimensional Hilbert spaces. Approach also provides a fruitful new perspective on arguments for the nonlocality and incompleteness of quantum theory. Specifically, **Nicholas Harrigan, Robert W. Spekkens** show that for models wherein the quantum state has the status of something real, the failure of locality can be established through an argument considerably more straightforward than Bell's theorem. The historical significance of this result becomes evident when one recognizes that the same reasoning is present in Einstein's preferred argument for incompleteness, which dates back to 1935. This fact suggests that Einstein was seeking not just any completion of quantum theory, but one wherein quantum states are solely representative of our knowledge. Hypothesis is supported by an analysis of Einstein's attempts to clarify his views on quantum theory and the circumstance of his otherwise puzzling abandonment of an even simpler argument for incompleteness from 1927. **February 2010, Volume 40, Issue 2, pp 125-157 Date: 09 Jan 2010 Einstein, Incompleteness, and the Epistemic View of Quantum States Nicholas Harrigan, Robert W. Spekkens** statistical model of the probabilistic description of physical reality has been studied by many authors. **Andrei Khrennikov posits** a contextual statistical model of the probabilistic description of physical reality. Here contexts (complexes of physical conditions) are considered as basic elements of reality. There is discussed the relation with QM. He proposes a realistic analogue of Bohr's principle of complementarity. In the opposite of the Bohr's principle, **Andrei Khrennikov** principle has no direct relation with mutual exclusivity for

observables. To distinguish his principle from the Bohr's principle and to give better characterization, he changes the terminology and speaks about supplementarity, instead of complementarity. Supplementarity is based on the interference of probabilities. It has quantitative expression trough a coefficient which can be easily calculated from experimental statistical data. There is need not appeal to the Hilbert space formalism and noncommutativity of operators representing observables. Moreover, in our model there exists a pair of supplementary observables which cannot be represented in the complex Hilbert space. There are discussed applications of the principle of supplementarity outside quantum physics. **October 2005, Volume 35, Issue 10, pp 1655-1693 The Principle of Supplementarity: A Contextual Probabilistic Viewpoint to Complementarity, the Interference of Probabilities and Incompatibility of Variables in Quantum Mechanics Andrei Khrennikov Foundations of Physics** Description of the vacuum in Yang-Mills theory and arrive at a physical interpretation of the pseudoparticle solution and the attendant violation of symmetries has been done by **R. Jackiw and C. Rebbi.** It is here vaccuum stability forms the bastion, pillar, post and stylobate of the mechanisms of yang mills Theory and action of a particle. The existence of topologically inequivalent classical gauge fields gives rise to a family of quantum mechanical vacua, parametrized by a CP-nonconserving angle. The requirement of **vacuum stability** against gauge transformations renders the vacua chirally noninvariant. DOI: http://dx.doi.org/10.1103/PhysRevLett.37.172 **Vacuum Periodicity in a Yang-Mills Quantum Theory Phys. Rev. Lett 37, 172 – Published 19 July 1976 R. Jackiw and C. Rebbi.** Three-dimensional Yang-Mills and gravity theories augmented by gauge-invariant mass terms are analyzed. These topologically nontrivial additions profoundly alter the particle content of the models and lead to (eb) quantization of a dimensionless mass-coupling-constant ratio. The vector field excitations become (=) massive, with (e&eb) spin 1 (rather than massless with spin 0), and the mass provides (eb) an infrared cutoff. The gravitation

acquires (e) mass, mediates (e&eb) finite-range interactions, and has (e) spin 2 (rather than being absent altogether); although its mass term is of (e)third derivative order, there are no ghosts or acausalities. DOI: http://dx.doi.org/10.1103/PhysRevLett.48.975 © 1982 The American Physical Society **Three-Dimensional Massive Gauge Theories Phys. Rev. Lett 48, 975 – Published 12 April 1982 S. Deser et al** All scattering amplitudes in (eb) the maximally supersymmetric N=4 super-Yang–Mills theory possess (e) a new, dual superconformal symmetry which extends the previously discovered dual conformal symmetry of MHV amplitudes. To reveal this property authors formulate the scattering amplitudes as (=) functions on (e&eb) the appropriate dual superspace. Rewritten in this form, all tree-level MHV and next-to-MHV amplitudes exhibit manifest dual superconformal symmetry. Authors propose a new, compact and Lorentz covariant formula for (e) the tree-level NMHV amplitudes for (e) arbitrary numbers and types of external particles. The dual superconformal symmetry is broken at loop level by infrared divergences. However, authors provide evidence that the dual conformal anomaly of (e) the MHV and NMHV superamplitudes is (=) the same and, therefore, their ratio is (=) dual conformally invariant. Authors show this explicitly for (e) the six-particle amplitudes at one loop. Authors conjecture that these properties hold for all, MHV and non-MHV, superamplitudes in N=4 SYM both at (eb) weak and at strong coupling. **Nuclear Physics B Volume 828, Issues 1–2, 21 March 2010, Pages 317–374 Dual superconformal symmetry of scattering amplitudes in N=4 super-Yang–Mills theory J.M. Drummonda DOI: 10.1016/j.nuclphysb.2009.11.022.** By comparison with numerical results in (eb) the maximal Abelian projection of lattice Yang–Mills theory, it is argued that the nonperturbative dynamics of (e) Yang Mills theory can be described by (e) a set of fields that take their values in (eb) the coset space SU(2)/U(1). The Yang–Mills connection is parameterized in (eb) a special way to separate (e&eb) the dependence on the coset field. The coset field is then regarded as (=) a collective variable and a method to obtain its effective

action is developed. It is argued that the physical excitations of (e) the effective action may be (=) knot solitons. A procedure to calculate the mass scale of knot solitons is discussed for (e) lattice gauge theories in the maximal Abelian projection. The approach is extended to (e&eb) the SU (N) Yang–Mills theory. A relation between the large N limit and the monopole dominance is pointed out. **Physics Letters B Volume 458, Issues 2–3, 8 July 1999, Pages 322–330 An effective action for monopoles and knot solitons in Yang–Mills theory Sergei V. Shabanov DOI: 10.1016/S0370-2693(99)00612-7.** It is here the quintessential formulation of consciousness occurs. Like a man standing on the threshold of infinity trying to ponder what lies beyond the veil, which separates the seen from unseen a true seeker tries to lift the individual general ledger to the cosmic general ledger. Like in Darwinian evolution, mind also has to evolve from individual consciousness to cosmic consciousness. In Kaishmiri Shaivism one withdraws from the finite to the infinite, but one also goes on an outward journey from the infinite to the finite, because both the finite and the infinite have an intimate connection. The finite is not seen as unreal, but as a symbol of the infinite. There is no real distinction between them. Those two movements constitute Spanda, a key concept in Kaishmiri Shaivism. Spanda is the pulsation of the Absolute in different phases of being. There are no opposites like subject and object, unity or duality, absolute or relative. They are just different phases of the universal vibration of the Absolute. The goal is to realize or be (=) at once infinite and finite. One does not turn away from appearances (like in Advaita Vedanta), but one realizes that the Absolute manifests all things. Spanda, the eternal pulsation of the Absolute, oscillates between a passion to create and dispassion from the created. Through it the Absolute transform itself into all things and then returns back into the emptiness of its undifferentiated nature. In Kaishmiri Shaivism the Absolute is seen as pure consciousness (=being). The Absolute is an eternal all-pervasive principle, the highest reality, the nature of all entities eternally and blissfully at rest within its own nature. The Absolute is the nature of the Self (and

thus of us all). The Absolute is divine; it is Shiva, the Lord of the Universe. It is full of (e) conscious activity through (e&eb)which it generates the universe, and reabsorbs it into (e&eb) itself at the end of each cycle of creation. Thus we speak of monism, as everything resides within (eb) this one absolute consciousness. It sustains (eb) all things, it embraces (e) all things, and it pervades (e&eb) all things. All things are appearances within (eb) the absolute consciousness, but nevertheless real (in contrast to (e) Advaita Vedanta where (e) appearances are seen as unreal or illusionary). All things appear external (out there, outside ourselves), but they do not have (e) a being on their own. They do not exist as (=) separate entities on their own. Everything is contained within (eb) consciousness. What we see as objects are (=) manifestations of consciousness. The event which constitutes (eb) the universe is always internal events happening within (e&eb) consciousness because (e) their essential nature is consciousness itself. If a physical object were totally material, and independent or external to consciousness, it could never be experienced. The universe and (e&eb) consciousness are two aspects of a whole. The universe is (=) an attribute of consciousness which bears consciousness as (=) its substance. Doctrine of Vibration: S.G.Dyczkowski what you "see" is (=) not always what you "get." Many people mistakenly take their own visions literally without (e) "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in (eb) the physical sense but in a psychological sense. The archetypes of (e) the collective unconscious are arrayed behind (e) the scenes of current worldwide conditions, of (e) crisis and confusion. They mirror (e&eb) our own states back at us, whether we perceive them as (=) such or interpret them plausibly or not (Note the essay has been written in model able form and models are given in various papers.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4\,a_9 \square - b_9 \square 0$$

And

$$b_9^{2} \square 4\,a_9 c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9\, in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\, X_9$ ,

$Let\ Y_9 = \ y_9\, X_9$And try the substitution $A_9 = X_9/(2 + A_9\ y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + \ n_9 y_9)$ we have

$$W_9^{2}\, y_9 \left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^{2}\left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9\, X_9 \left[\frac{X_9 (1 + n_9 y_9)}{M_9^{2}(2 + n_9 y_9)^{2}} + 2\, y_9 \left(n_9 - 1\right) + \frac{X_9^{2}(9 + y_9)}{M_9^{2}(2 + n_9 y_9)} = 0.\ (A2.12)\right.$$

Since $n_9$ and $A_9$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9 = (2 + n_9 y_9)$, *the condition*(A2.99?)*can be written as*

$[\frac{X_9(E_9-9)}{E_9^2}]^2[\frac{X_9^2}{M_9^2} - 4\, y_9\ E_9^2 \frac{X_9}{M_9^2} + 4y_9 (n_9 - 2)E_9^3]$ .

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9 / W_9^2)$ can be expressed as

$2\, E_9 [y_9 E_9 + (y_9\ E_9 (9 + y_9)(E_9 - n_9))^{9/2}] > \frac{X_9}{M_9^2} > 2E_9 [y_9 E_9 - y_9 E_9 (9 + y_9)(E_9 - n_9))^{9/2}]$

Concept Of Being From Shaivite Point Of View: In His Famous Work Entitled On The Conservation Of Force (1847), Hermann Von Helmholtz Had Divided The Totality Of The Energy In The Universe Between Potential Energy And Kinetic Energy And Affirmed The Reciprocal Convertibility Of The Two. In 1852, William Thomson Pointed Out That There Exists A Sub-Ensemble Within Kinetic Energy, Heat, Which, Once It Has Been Generated, Is No Longer Entirely Convertible Into Potential Energy-Or Into Any Other Form Of Kinetic Energy. Considering That The (Partial) Reconversion Of Heat Into Labor Is Possible Only In Situations That Present A Disparity In Temperature, And That Heat Tends To Pass From Warmer To Cooler Bodies By Spreading On An Even Temperature Level Through Space, Thomson Concluded That The Universe Tends Towards A Final State Where Any Energetic Transformations, Every Movement And Every Form Of Life Will Cease: The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO: The Ultimate Nature Of Anything Cannot Change. Change Of Any Sort Is Merely Apparent (Vivarta). The World Of Change And Becoming Is A False Super- Imposition (Adhyaropa, Adhyasa) On The Absolute. It Appears So On The Screen Of Individual, Collective Consciousness, Albeit One Cannot Be So Sure Of Cosmic, For There Are Many Reasons That Cosmic Transactions Also Had To Be Subjected To Change

By ParaBrahaman. In Cosmic Terms, The Mistake (Bhranti) Consists Of The Supposition That The Real Brahman Is The Unreal Universe And The Unreal Universe Is The Real Brahman. In Microcosmic Terms, It Is The Mistake Of Falsely Conceiving The Body, Mind Or Even One's Personality To Be The Self. Note We Have Taken Self As Witness Consciousness. Notwithstanding There Is Nothing Wrong In The Conception Of Self As A Dynamical Presupposition, In That Individual Consciousness And Collective Consciousness And Even Cosmic Consciousness Evolutes In To ParaBrahaman. In The Same Way As The Image Of A Snake Is Falsely Superimposed On A Rope, Similarly The Universe Is Falsely Projected Onto The Real Substratum, The Brahman. Ignorance Is Not Merely A Personal Lack Of Knowledge, But A Cosmic Principle. As Such It Is Called "Maya," The Indefinable Factor (Anirvacaniya) That Brings This Mistake In Identity About. The Reality Status Of This Cosmic Illusion Is Also Indefinable: On The One Hand It Is Not Brahman, The Sole Reality; On The Other Hand It Is Not Absolutely Non-Existent Like A Hare's Horn Or The Son Of A Barren Woman. Note Here The Vedantin Avoids Talking About Zero Which We Have Avidly And Assiduously Stressed About. Nothing Is The Factor That Exists. Also See the Remarks about Satre's Being And Nothingness. Brahman Is The Source Of World Appearances Only In The Sense Of Being Their Unconditioned Ground Or Essential Nature. The Universe Is False Not Because It Has No Nature Of Its Own But Because It Does Have One. The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski. Concept Of Being From Shaivite Point Of View: It Is The 'I Am The Body/Mind' Belief That Gives Rise To The 'I Am Not The World' Belief. These Two Beliefs Are (=) Co-Created."Consciousness Projects The Appearance Of The Mind, Body And World By Taking (Eb) The Shape Of Thinking, Sensing And Perceiving."“Attention Is Consciousness With An Object. When The Object Vanishes, Attention Simply Remains What It Always Is Consciousness."“There Is No Purpose To Meditation. The Purpose Is Already Accomplished."“Everything

That Is Experienced Is Experienced By, Through, In And As Consciousness.”“The Seen Cannot Be Separated From Seeing And Seeing Cannot Be Separated From Consciousness.” “**The Reality Of Any Experience Is Not Hidden In The Appearance; It Is Expressed By The Appearance The Transparency Of Things**' By Rupert Spira Book Review By Dennis Waite The Following Is A Review Of Rupert's Book: 'The Transparency Of Things: Contemplating The Nature Of Experience’. There Is An Essay From The Book Here And The Essay On Cézanne May Be Read At The Awakened Eye Website Nature Rejects The Naiveté That Seeks Absolute Truth. We Are Beginning To Realize, Individually And Culturally, That "Realities" Are All Human Constructions. The Task Becomes One Of "Catching Ourselves In The Act" Of Creating Our Own "Reality" From The Flow Of Events. Human Truth Is Always An Engagement Of Mind With Experience. The Challenge Of The Therapist In These Times Of Chaotic Change Is To Validate The Concept That We Don't Have To Fear The Collapse Of What We Think We Are. Strange Attractors: Transference, Holography, And An **Archetype Burke, J. (2003). Strange Attractors: Transference, Holography, And An Archetype (Doctoral Dissertation, Pacifica Graduate Institute, 2003).** When we get to the topic and understanding of the Transcendental Logic, the brilliant and baffling Transcendental Deduction (TD) Recall the two movements just discussed, the one from experience to its conditions and the one from the forms of valid inference to the concepts that we must use in all judging (the Categories). This duality led Kant to his famous question of right (quid juris) (A84=B116): with what right do we apply the Categories, which are not acquired from experience, to the contents of experience? (A85=B117). Kant's problem here is not as arcane as it might seem. It reflects an important question: How is it that the world as we experience it conforms to our logic? In briefest form, Kant thought that the trick to showing how it is possible for the Categories to apply to experience is to show that it is necessary that they apply (A97).TD has two sides, though Kant never treats them separately John Marlado,

accounts, attributes, and ascribes the philosophy of Nishida KitarÅ is lucid, insightful, and deeply informative despite the highly questionable argument that gives the book its structure. In comparison with other, technically more accurate accounts of Nishida's philosophy, it is broader in scope than most but more single-minded in its approach. The analysis is for the most part limited to four books that have been translated into English, out of approximately 16 volumes of philosophical essays in Japanese, but it aims to capture the contours and most important details of Nishida's entire philosophy. To appreciate the achievement of this account, one would need only to read, in English translation or in the Japanese original, any of the volumes with their streams of repetitive, puzzling expressions and circuitous, often dead- end argumentation. Wilkinson's elucidation of details in the passages he samples makes it easier to understand why Nishida is considered Japan's greatest academic philosopher. A critical look at the frame of the book and some of its own claims reveals where, to my mind, it has gone astray. The argument that frames this book is repeated throughout: Nishida's philosophy is an attempt to articulate, in a characteristically Western philosophical manner, the Zen experience he had as a young man. Nishida first made the attempt by using the conceptual frameworks of Western philosophers, and after they failed him he developed his own, largely by recasting Buddhist insights. In the end, a philosophical articulation of the Zen worldview in Western terms or even in Eastern terms recast by conceptual, Western philosophical methods approved impossible. What Nishida attempted is incommensurable with Western philosophy. While the author first offers these statements as hypotheses (p. 2), one premise of his argument that there is such a thing as "the Zen world-view" or "the Zen experience “remains unquestioned; and another that Nishida's pivotal experience was the Zen experienceâis asserted as "beyond question" (p. 151; see also pp. 28, 48, 10, 114).Wilkinson's presentation of Zen in the first chapter is drawn largely from D. T. Suzuki, with help from Suzuki's portage Abe Masao; some comments on time, birth, and death in; and a few other sources. The

result is a caricature that imagines Zen as an invariable, unitary experience. In fact, Dagan's teachings on practice as the manifestation of enlightenment do not sit well with D. T. Suzuki's emphasis on a satori experience, and Suzuki does not enlighten us about what Zen in its various practices has historically been. The author would have benefited from acquaintance with other presentations by Zen teachers, such as that of ShunryÅ Suzuki'sZen Mind Beginner's Mind, with historical accounts of Zen traditions, and with contrasts between DÅgen's view of time and Sam Â· sÄraand that of other Buddhist philosophers. A little research in contemporary Zen scholarship would have shown how implausible is the talk of "nirvana or awareness of mu" (p. 24), "the basic Zen position [that] never varies" (p. 3), "the Zen vision" (pp. 5, 7), "Zen epistemology" , "Zen experience" (p. 59), "the conception of time involved in Zen" (p. 15), and "the fundamental Zen assertion that there is no absolute time" (p. 16). Even if DÅgen's notion that time and being are nondual (p. 16) epitomized most of Zen, Nishida's early notion of "a transcendent, unchanging reality apart from time" (p. 51) would contrast sharply with it. Equally questionable is Wilkinson's contention that Nishida's "pure experience" is equivalent to Zen satori. The notion of pure experience Nishida developed in his first major work may be pivotal for his entire philosophy, but he never claims or implies that any part of his philosophy presupposes an exceptional experience like satori that grounds his convictions and renders them unverifiable to those who have not had the experience, as Wilkinson intimates (pp. 55, 159, and 161). Nishida presents his notions of pure experience and intellectual intuition, the grasp of its unity, and later... Perennialism and Constructivism is studied by Randolph T Dible : experiencer or the observer and the transcendental subjectivity itself Pure objectivity is essentially a deconstruction, apophatic and negativa The encountering of the experiencer or observer—transcendental subjectivity itself—at the foundation of the world leads inevitably to the recognition of pure objectivity as ultimate reality (which can be taken as its ultimate deconstruction, analogous to the apophatic or via

negativa), from which objects derive their value, weight, significance, meaning or objectivity. In this way, pure objectivity can be seen as the supra-self-evident Axiological Axiom, so to speak, even Unconditional Love, in romantic terms. This axiology (value theory) has a structure inverse to the relationship between transcendental subjectivity as the radical unity of pure self-reference and on the other hand, the world of forms, as mere traces (representations, indications) of the unique, original "first distinction" Spencer-Brown speaks of at the foundation of his calculus. That is, all forms (i.e., distinctions, differences) would reduce to being the first distinction, also known as the marked state, which can be called penultimate reality (pure self-reference or transcendental subjectivity: the Spirit which animates us), except that forms are complimentary to their content, which is their objectivity or value, which would reduce to the unmarked state or ultimate reality. It is the incongruity of form (thoughts; Whitehead's "negative prehensions") and value (feelings; Whitehead's "positive prehensions," or my notion of objectivity, meaning and qualia; in short, the non-formal aspects of experience) that holds forms open and keeps them from absolute reduction. This accounts for the brute, concrete persistence of the "functional illusion"-- to use a term from Dzogchen Buddhism-- of the world. Thus this system has an axiology of metaphysical objectivity grounded on the ideal of pure objectivity as the source of all value, meaning and significance, itself the very fecundity of profundity, which is the motive of drawing the distinction in the first place **(Wikipedia, Kant's writings and Logic Of Sense)** Nano rings were created by accident while intending to make quantum dots. They have interesting optical properties associated with excitons and the Aharonov–Bohm effect. Application of these rings used as light capacitors or a buffer includes photonic computing and communications technology. Analysis and measurement of geometric phases in mesoscopic rings is ongoing. Several experiments, including some reported in 2012, show Aharonov-Bohm oscillations in charge density wave (CDW) current versus magnetic flux, of dominant period h/2e through CDW rings up to 85 **μm** in

circumference above 77 K. This behavior is similar to that of the superconducting quantum interference devices (seeSQUID). **Aharonov–Bohm nano rings (Wikipedia) Stability analysis of nano rings would help progressively accentuate the research about the materials division that uses nano rings.** From the profound contemplation and wisdom of the Buddha and Mahabodhisattvas, we know there is no such reality. Instead, egolessness (non-self) is the only path to understand the reality of the deluded life. All existences are subject to the law of causes and conditions. These include the smallest particles, the relationship between the particles, the planets, and the relationship between them, up to and including the whole universe! From the smallest particles to the biggest matter, there exists no absolute independent identity. Egolessness (non-self) implies the void characteristics of all existence. Egolessness (non-self) signifies (eb) the non-existence of permanent identity for self and existence (dharma). Sunyata stresses (eb) the voidness characteristic of (e) self and existence (dharma). Sunyata and Egolessness possess (e) similar attributes. As we have discussed before, we can observe the profound significance of Sunyata from (e) the perspective of inter-dependent relationships. Considering dharma-nature and (e&eb) the condition of nirvana, all existences are immaterial and of (e) a void-nature. Then we see each existence as (=) independent of each other. But then we cannot find any material that does exist independent of (e) everything else. So Egolessness also implies (eb) void-nature! From the observation of all existences, we can infer (eb) the theory of nirvana and the complete cessation of (e) all phenomena. From the viewpoint of phenomena, all existences are so different from (e) each other, that they may contradict (e&eb) each other. They are (=) so chaotic. In reality, their existence is (=) illusionary and arises from (e) conditional causation. They seem to exist on one hand, and yet do not (e) exist on the other. They seem to be united, but yet they are so different to (e) one another. They seem to exist and yet they do cease! Ultimately everything will return to (e) harmony and complete calmness. This is the nature of all existence. It is the

final resting place for all. If we can understand this reality and remove our illusions, we can find this state of harmony and complete calmness. **Teachings in Chinese Buddhism (6) sunyata (emptiness) in the Mahayana context (Wikipedia) Many models have been given concatenating Buddhist doctrines and Vaccum energy. These models essentially throw light on the essence of self(witness consciousness) and the cosmic general ledger (cosmic consciousness) while the imbalance in the gratification deprivation complex always remain unbalanced when it comes to individual general ledger , advocating the evolutionary theory of individual consciousness to cosmic consciousness. It is to be noted that it is not the question of knowing all the transactions but the wisdom of what lies behind the debit- Credit, Debit-Debit, or Credit-Credit principle that is most important for assimilation.** Look at this simple thesis to prove the point stated in the foregoing. Thesis: 'The World Has A Beginning In Time, And In Space It Is Also Enclosed In Boundaries.' Proof: 'For If One Assumes That The World Has No Beginning In Time, Then Up To Every Given Moment In Time An Eternity Is Elapsed, And Hence An Infinite Series Of States Of Things In The World, Each Following Another, Has Passed Away. But Now The Infinity Of A Series Consists Precisely In The Fact That It Can Never Be Completed Through A Successive Synthesis. Therefore An Infinitely Elapsed World-Series Is Impossible, So A Beginning Of The World Is A Necessary Condition Of Its Existence, Which Was The First Point To Be Proved (Kant, Critique Of Pure Reason, Tr. By R. Guyer And A. W. Wood (Cambridge: Cambridge University Press, 1998): B 454, P. 470). As Quoted In The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO"The Cyclical Hypothesis, So Heavily Criticized By Nietzsche (VP II 325 And 334), Arises In This Way." In Fact, Nietzsche Was Not Criticizing The Cyclical Hypothesis But Only The Particular Form Of That Hypothesis Presented In Vogt's Work. All Of Nietzsche's Texts Without Exception Speak Of The Eternal Return As The Repetitio Of The Same Events Within A Cycle Which Repeats Itself

Eternally. If Deleuze's Interpretation Holds That The Eternal Return Is Not A Circle, Then What Is It? A Wheel Moving Centrifugally, Operating A "Creative Selection," "Nietzsche's Secret Is That The Eternal Return Is Selective" Says Deleuze: The Eternal Return Produces Becoming-Active. It Is Sufficient To Relate The Will To Nothingness To The Eternal Return In Order To Realize That Reactive Forces Do Not Return. However Far They Go, However Deep The Becoming-Reactive Of Forces, Reactive Forces Will Not Return. The Small, Petty, Reactive Man Will Not Return. Affirmation Alone Returns, This That Can Be Affirmed Alone Returns, Joy Alone Returns. Everything That Can Be Denied, Everything That Is Negation, Is Expelled Due To The Very Movement Of The Eternal Return. We Were Entitled To Dread That The Combinations Of Nihilism And Reactivity Would Eternally Return Too. The Eternal Return Must Be Compared To A Wheel; Yet, The Movement Of The Wheel Is Endowed With Centrifugal Powers That Drive Away The Entire Negative. Because Being Imposes Itself On Becoming, It Expels From Itself Everything That Contradicts Affirmation, All Forms Of Nihilism And Reactivity: Bad Conscience, Ressentiment..., We Shall Witness Them Only Once. [...] The Eternal Return Is The Repetition, But The Repetition That Selects, The Repetition That Saves. Here Is The Marvelous Secret Of A Selective And Liberating Repetition. There Is No Need To Remind The Reader That Neither The Image Of A Centrifugal Movement Nor The Concept Of A Negativity-Rejecting Repetition Appears Anywhere In Nietzsche's Writings, And Indeed Deleuze Does Not Refer To Any Text In Support Of This Interpretation. Further, One Could Highlight That Nietzsche Never Formulates The Opposition Between Active And Reactive Forces, Which Constitutes The Broader Framework Of Deleuze's Interpretation. For Some Years, Marco Brusotti Has Called Attention To The Fact That Deleuze Introduced A Dualism That Does Not Exist In Nietzsche's Writings. To Be Sure, The German Philosopher Describes A Certain Number Of "Reactive" Phenomena (For Example, In The Second Essay Of The Genealogy Of Morality, § 11, He Talks About

"Reactive Affects" [Reaktive Affekte], "Reactive Feelings" [Reaktive Gefühlen], Reactive Men [Reaktive Menschen]); But These Are Nonetheless The Result Of Complex Ensembles Of Configurations Of Centers Of Forces That Remain In Themselves Active. Neither The Word Nor The Concept Of "Reactive Forces" Ever Appears In Nietzsche's Philosophy. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO Some Parts Are Deleted Due To Spatial Constraints. In the midst of thoughts about the eternal return we find at least two other thematic axes On the one hand, the view of the **world as a constant flux of forces without any goal, law, or rules of becoming. A chaos sive natura de-divinized and de-anthropomorphized which constitutes the "ontological substratum**" of the whole of Nietzsche's reflexions On the other hand, an ensemble of fragments of an anthropologico-sociological character, designing a path of liberation leading to the creation of superior individuals by way of a profound transformation of their instinctual structure. This transformation must be achieved by a practice of solitude and internal struggle towards the liberation from the ancient representations of the world and from the incorporated herd values. For an analysis of these thematic perspectives, see Paolo D'Iorio, La linea e il circolo. Cosmologia e filosofia dell'eterno ritorno in Nietzsche (Genova: Pantograf, 1995): 233-322. With most of the parameters quantified, with assumptions there must be picturesque view of the self and non self, space and time, science and philosophy with the concatenational consummations. After The Chapter "On Redemption," Where Zarathustra Dares Not Expose His Doctrine, The Eternal Return Begins To Be Enunciated In Part Three Of The Work. In The First Place, It Is The Dwarf Who Formulates It In The Chapter "On The Vision And The Riddle." Facing The "Gate Of The Instant" Which Symbolizes The Two Infinities That Stretch Towards The Past And The Future, The Dwarf Whispers: "All Truth Is Crooked, Time Itself Is A Circle" The Dwarf Represents The Spirit Of Gravity, And He Embodies The Herd Morality, "The Belittling Virtue" Which Is The Title Of Another Chapter From Part III. The Dwarf Can

Endure The Eternal Return Without Great Difficulties Because He Has No Aspirations; Unlike Zarathustra He Does Not Wish To Climb The Mountains That Symbolize Elevation And Solitude. In Two Unpublished Notes, From The Summer And The Fall Of 1883, Nietzsche Writes: The Doctrine Is At First Favored By The Rabble, Before It Gets To The Superior Men. The Doctrine Of Recurrence Will First Smile To The Rabble, Which Is Cold And Without Any Strong Internal Need. It Is The Most Ordinary Of Life Instincts, Which Gives Its Agreement First. Hence, The Content Of The Doctrine Is The Same, But Whereas The Dwarf Can Endure It (Because He Interprets It According To The Pessimistic Tradition For Which "Nothing Is New Under The Sun"), Zarathustra, Who Is The "Advocate Of Life" Regards The Eternal Return As The Strongest Objection To Existence, And As The Rest Of The Dream Suggests, He Does Not Yet Succeed In Accepting It After The Vision At The Gate Of The Instant, The Chapter Is Brought To An End By The Enigma Of The Shepherd. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO .How does the memory affect the GD complex or for that matter the gratification deprivation of an individual or a collective identity? Non-Markovian local-in-time master equations give a relatively simple way to describe the dynamics of open quantum systems with memory effects. Despite their simple form, there are still many misunderstandings related to the physical applicability and interpretation of these equations. Here, we clarify these issues both in the cases of quantum and classical master equations. Further introduction of the concept of a classical non-Markov chain signified through negative jump rates in the chain configuration. **E-M Laine et al 2012 J. Phys. B: At. Mol. Opt. Phys. 45 154004 doi:10.1088/0953-4075/45/15/154004 Local-in-time master equations with memory effects: applicability and interpretation.** Supposing That There Were Indeed An "Energy Of Contraction" Constant In All Centers Of Force Of The Universe, It Remains To Be Explained Where Any Difference Would Ever Originate. It Would Be Necessary For The Whole To Dissolve Into An Infinite Number Of

Perfectly Identical Existential Rings And Spheres, And We Would Therefore Behold Innumerable And Perfectly Identical Worlds COEXISTING [Nietzsche Underlines This Word Twice] Alongside Each Other. Is It Necessary For Me To Admit This? Is It Necessary To Posit An Eternal Coexistence On Top Of The Eternal Succession Of Identical Worlds? The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. Science is usually portrayed as dehumanizing. Brave New World epitomizes this fear. "The more we understand the world, the more it seems completely pointless" (Steven Weinberg). Certainly science can seem chilling when conceived in the abstract as a metaphysical world-picture. We may seem to find ourselves living in a universe with all the human meaning stripped out: Participants in a soulless dance of molecules, or harmonics of pointlessly waggling superstrings and their braneworld cousins. Nature seems loveless and indifferent to our lives. What rights have we to be happy? **(Brave New World: Aldous Huxley).** Religion can afford to be dogmatic but science cannot. Look at this passage which counters both scientists and philosophers alike: However, Hartmann Objects That The Regressive Movement Postulated By Schopenhauer Is Possible Only In Thought: It Remains Nothing More Than An "Ideal Postulate" With No Real Object And Which "Does Not Teach Us Anything About The Real Process Of The World That Unfolds In A Movement Contrary To This Backwards Movement Of Thought" (Hartmann, Philosophie Des Unbewussten, Third Edition (1871): 772). Hartmann Affirms That If Unlike Schopenhauer One Admits The Reality Of Time And Of The World Process, One Must Also Admit That The Process Must Be Limited In The Past And Therefore That There Must Be An Absolute Beginning. In Hartmann's Mind, Failure To Do So Would Result In Positing The Contradictory Concept Of An Accomplished Infinity: "The Infinity That From The Point Of View Of Regressive Thinking, Remains An Ideal Postulate, Which No Reality May Correspond To, Must, For The World, Whose Process Is, On The Contrary, A Progressive Movement, Open Up To A Determinate Result; And Here The Contradiction Comes To Light" (Hartmann

(1871): 772). What Really "Comes To Light" In This Passage Is The Fact That Hartmann Does Not Provide A Demonstration But A Petitio Principii. Indeed, The Concept Of The World Process Analytically Contains The Concept Of A Beginning Of The World. In All Rigors, It Is Therefore Impossible To Demonstrate These Concepts With Reference To Each Other. Secondly, Hartmann's View That One Is Bound To Accept The Reality Of The World Process Even If One Rejects The Ideality Of Schopenhauer's Time Is Mistaken. Hartmann Believes That If Time Is Real There Must Be A World Process With Both An Absolute Beginning And An Absolute End. Without Any Justification, Hartmann Jumps From Schopenhauer's Negated Time To Oriented Time. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. In the last years several theoretical papers discussed if time can be an emergent property deriving from ) quantum correlations. Here, to provide an insight into how this phenomenon can occur, authors present an experiment that illustrates Page and Wootters' mechanism of "static" time, and Gambini et al. subsequent refinements. A static, entangled state between a clock system and the rest of the universe is perceived as evolving by internal observers that test the correlations between the two subsystems. They implement this mechanism using an entangled state of the polarization of two photons, one of which is used as a clock to gauge the evolution of the second: an "internal" observer that becomes correlated with the clock photon sees the other system evolve, while an "external" observer that only observes global properties of the two photons can prove it is static.Phys. Rev. A 89, 052122 (2014)DOI:10.1103/PhysRevA.89.052122 arXiv: 1310.4691 [quant-ph] **Time from quantum entanglement: an experimental illustration Ekaterina Moreva.** Time is an entanglement .When the new ideas of quantum mechanics spread through science like wildfire in the first half of the 20th century, one of the first things physicists did was to apply them to gravity and general relativity. The results were not pretty. It immediately became clear that these two foundations of modern physics were entirely

incompatible. When physicists attempted to meld the approaches, the resulting equations were bedeviled with infinities making it impossible to make sense of the results. Then in the mid-1960s, there was a breakthrough. The physicists John Wheeler and Bryce DeWitt successfully combined the previously incompatible ideas in a key result that has since become known as the Wheeler-DeWitt equation. This is important because it avoids the troublesome infinites—a huge advance. But it didn't take physicists long to realise that while the Wheeler-DeWitt equation solved one significant problem, it introduced another. The new problem was that time played no role in this equation. In effect, it says that nothing ever happens in the universe, a prediction that is clearly at odds with the observational evidence. This conundrum, which physicists call 'the problem of time', has proved to be thorn in flesh of modern physicists, who have tried to ignore it but with little success. Then in 1983, the theorists Don Page and William Wootters came up with a novel solution based on the quantum phenomenon of entanglement. This is the exotic property in which two quantum particles share the same existence, even though they are physically separated. Entanglement is a deep and powerful link and Page and Wootters showed how it can be used (e) to measure time. Their idea was that the way pair of entangled particles evolves is a kind of clock that can be used to measure change. But the results depend on how the observation is made. One way to do this is to compare the change in the entangled particles with an external clock that is entirely independent of the universe. This is equivalent to god-like observer outside the universe measuring the evolution of the particles using an external clock. In this case, Page and Wootters showed that the particles would appear entirely unchanging—that time would not exist in this scenario. But there is another way to do it that gives a different result. This is for an observer inside the universe to compare the evolution of the particles with the rest of the universe. In this case, the internal observer would see a change and this difference in the evolution of entangled particles compared with everything else is an important a measure of time. This is

an elegant and powerful idea. It suggests that time is an emergent phenomenon that comes about because of the nature of entanglement. And it exists only for observers inside the universe. **Time is an entanglement** (Wikipedia). Danish religious philosopher Soren Kierkegaard carried out a systematic critique of the pretensions of reason and an abstract rationalism which he believed that the modern age was nurturing. Condemning reflection as a "danger" that ensnares people in logical delays and machinations, Kierkegaard compared it to a prison. Reflection is for him a form of captivity, a bondage which "can only be broken by [passionate] religious inwardness" (1978: 81). **Reflection seduces individuals into thinking its possibilities are "much more magnificent than a paltry decision"** (1978: 82). It leads them to act **"on principle,"** to dwell on the deliberation of the context of their actions and the calculation of their worth or outcome. Kierkegaard argues that this drives away feeling, inspiration, and spontaneity, all of which are crucial for true inner being and a vital relation to God. For Kierkegaard, as Nietzsche would later agree, genuine inner being (and culture) is characterized by the tautness and tension of the soul which characterizes passionate existence. But the "coiled springs of life relationships ... lose their resilience" in reflection (1978: 78) and "everything becomes meaningless externality, devoid of [internal] character" (1978: 62). Kierkegaard thus contributes to the development of an irrationalist tradition that has echoes in some later postmodern thought. Kierkegaard might well have agreed with his contemporary Fyodor Dostoyevsky, who wrote: "An intelligent [reflective] man cannot seriously become anything ... excessive consciousness is a disease" (1974: 3, 5). In an age overtaken by rules and regulations, genuine action -- which Kierkegaard assumes to be subjective and spontaneous -- is frustrated at every turn. Complaining that we are too "sober and serious" (1978: 71) even at banquets, Kierkegaard bemoans the fact that even suicides are premeditated (1978: 68)! "That a person stands or falls on his actions is becoming obsolete; instead, everybody sits around and does a brilliant job of bungling through with the aid of some reflection

and also by declaring that they all know very well what has to be done" (1978: 73). Thus, it is passion, not reflection, that guarantees "a decent modesty between man and man [and] prevents crude aggressiveness" (1978: 62). "Take away the passion and the propriety also disappears" (1978: 64). The ambiguity in the word "passion" may cause some confusion here. To say that the age and its individuals are "passionless" is not to say there are no emotions whatsoever, but rather that there is no true spiritual inwardness and depth, no intensively motivated action and commitment. It suggests that passion exists only in a simulated, pseudo-form, "the rebirth of passion" through "talkativeness" (1978: 64). "Chattering" for Kierkegaard gets in the way of "essential speaking" and merely "reflects" inconsequential events (1978: 89-99). Hence, in "the present age," emotions -- which in fact are all too pronounced -- have been transformed into negative forces. Anticipating Nietzsche's genealogy of the "slave revolt" in morality, Kierkegaard claims that the "enthusiasm" of the prior age of Revolution, a "positively unifying principle," has become a vicious "envy," a "negatively unifying principle" (1978: 81), a leveling force in its own right insofar as those lacking in talent and resources want to tear down those who have them. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS By Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html.** In November 2011, the United States government argued before the US Supreme Court that it wants to continue utilizing GPS tracking of individuals without first seeking a warrant. In response, Justice Stephen Breyer questioned what this means for a democratic society by referencing Nineteen Eighty-Four. Justice Breyer asked, "If you win this case, then there is nothing to prevent the police or the government from monitoring 24 hours a day the public movement of every citizen of the United States. So if you win, you suddenly produce what sounds like 1984... "[65]. What is to be said in unmistakable terms is that when it comes to national security and innocent citizens there is question of compromise even it be an infringement and encroachment of so called

privacy which has been made a word to subterfuge the activities that are far said to be fair. Quantum theory has provoked intense discussions about its interpretation since its pioneer days. One of the few scientists who have been continuously engaged in this development from both physical and philosophical perspectives is Carl Friedrich von Weizsäcker. The questions he posed were and are inspiring for many, including the authors of this contribution. Weizsäcker developed Bohr's view of quantum theory as a theory of knowledge

## ABSOLUTE STABILITY

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range (0 W 9 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9 - 9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2 - 2y_9\frac{(E_9 + n_9^2 y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right) - 4y_9(E_9 - n_9)$$ 0

And has the solution

$$\frac{X_9}{M_9^2} > \frac{E_9}{(E_9 - 9)}[y_9(E_9 + n_9^2 y_9) + [y_9^2 (E_9 + n_9^2 y_9)^2 + 4y_9(E_9 - 9)(E_9 - n_9)]^{9/2}]$$

Given $E_9 > n_9$this lower limit to $X_9 / M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9 / M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9 < n_9$ $(ie., 2A_9 M_9^2 < X_9 - Y_9$)then (A2.9)and (A2.8)show that $Q_9$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2} > 2E_9[y_9 E_9 - (y_9 E_9 (9 + y_9)(E_9 - n_9)]^{9/2}]$$

And

$$E_9 > n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2/2M_9 (X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$and $X_9$and $Y_9$by $W_9^2$ the dependence upon *M* disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9$ *and* separate (A3.1)into real and imaginary parts

$$Z_9^2 + (M_9 + V_9) Z_9 \sin(Z_9 T_9) - (U_9 M_9 + Y_9) + (M_9(U_9 + V_9) - X_9)$$

$$\mathrm{X}\cos(Z_9 T_9) - 2M_9 V_9 \cos^2(Z_9 T_9) + M_9 V_9 = 0$$

$$(M_9 + U_9) Z_9 - (M_9 + V_9) Z_9 \cos(Z_9 T_9) + (M_9(U_9 + V_9) - X)$$

$$\mathrm{X}\sin((Z_9 T_9) - M_9 V_9 \sin(Z_9 T_9) = 0$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 = 1 - \cos Z_9 T_9$and $A_9 = U_9 - V_9$we obtain

$$4 M_9^2 V_9^2 W_9^4 + 2 M_9 V_9 W_9^3[(M_9 - V_9)^2 + A_9(M_9 + V_9) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2[M_9 - V_9)^2(2 A_9 M_9 + Y_9 - X_9) + (2A_9 M_9 - X_9)^2 - A_9(M_9 - V_9)(2A_9 M_9 - X_9)$$

$$+ M_9 V_9 (A_9^2 + 4(X_9 + Y_9))]$$

$$+ W_9[A_9^2 + 2X_9)(2 A_9 M_9 - X_9) + 2A_9 Y_9(M_9 + V_9) - M_9 A_9^3] + A_9^2(X_9 M_9 - Y_9) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9(2, U_9 V_9) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting cos $(Z_9 T_9) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MDDULE NUMBERED SEVEN**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability

criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

## Morphological And Epistemological And Syllogistic Models For Which Nunny's Model Could Be Applied And Has Been Applied:

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production**

**Charge Asymmetry, $\Delta Y \equiv Yt - Yt^{-}$, Reflects The Asymmetry In $Tt^{-}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles

ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic

quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation

between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to

nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of

Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first

impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of

another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease. Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation

of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces

always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the

thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of

thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity

and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI

Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs” Outcomes—

And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities

Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

## DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t-T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\big(R(t)\big)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and haimovici model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

**TIME DELAY**

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Limitations of Predator prey analysis in immense. When we realize that world itself is Debit-Credit world and sometime credit-Credit world (See an earlier paper entitled the same nomeneclature.

Opposite limit of absolute stability. The result is a specific The Doctrine of Vibration awareness of an object of the form: 'here this is such.' While individual appearances do not lose their separate identity {svarupabheda) when they rest on a common basis, even so the particular object which appears according to its own characteristics (svalak$ana) is an individual reality in its own right. It is a different kind of appearance characterised by its association with the appearance of the specific location and time in which it is made manifest. The form of our experience is thus 'I now see this here'. But when we perceive each particular constituent appearance separately, each assumes a separate fixed function. Abhinava cites the following colourful example to illustrate how the various combinations of appearances account for the variety of experience: Thus even though the appearance of the beloved may manifest externally, it is as if far away in the absence of another appearance, namely, that of 'embracing'. So when the [appearing of the beloved] is associated with another appearance [namely that of 'far away'] the power (arthakriya) it formerly had of giving pleasure appears as its contrary. The form our experience assumes depends, not only on the nature of the object perceived, but also on personal factors entirely peculiar to ourselves. This theory explains this in two ways. In one sense, the object remains the same, but one or other of its constituent appearances

comes to the fore according to the inclinations of the perceiver. From another point of view, we can say that the perceived object is different for each perceiver according to the difference in the prominent appearance manifest to him. Abhinava, citing as an example a golden jar, illustrates how the same object appears differently to different perceivers according to the use they wish to make of it and to their state of mind: When a person who is depressed and feels that there is nothing [of value for him in the world] sees the jar, he merely perceives the appearance 'exists' [in the form of the awareness that] 'it is\ He is not conscious of any other [of its constituent appearances] at all. An individual who desires to fetch water [perceives] the appearance 'jar'. The man who simply wants something that can be taken somewhere and then brought back [perceives] the appearance 'thing'. The man who desires money [perceives] the appearance 'gold'. The man who desires a pleasing object [perceives] the appearance 'brightness' while he who wants something solid sees the appearance 'hardness'. These 'atomic events' or appearances emerge from the pure subject's consciousness and combine together to form a total event at each moment. Integral Monism of Kashmiri Saivism Daily life (vyavahara) goes on by virtue of this ever renewed flux of appearances. They are connected together and work towards a single unified experience because they appear within the field of consciousness of the universal subject. The aggregate of appearances arises in the [supreme] subject as do [sprouts in] a rice field. Even though each sprout germinates from its own seed, they are perceived as a collective whole. Appearances rest in this way within the universal subject. 'External- ity' is itself another appearance; it arises from a distinction between appearances and the individual subject. So, although all manifestation always occurs within the subject, it appears to be external due to the power of Maya which separates the individual subject from his object. This split must occur for daily life to be possible. Only externally manifest appearances can perform their functions; when they are merged within the subject and at one with him, they cannot do so. Daily life proceeds on the basis of the operation and withdrawal of the conditions necessary for fruitful action to be possible. Appearance in this sense represents the actualisation of a potential hidden in consciousness made possible by virtue of its dynamic, Spanda

nature which is both the flow from inner to outer and back as well as the power that impels it. The emergence from, and submergence into, pure consciousness of each individual appearance is a particular pulsation (visesaspanda) of differentiated awareness. Together these individual pulsations constitute the universal pulse (samanyaspandd) of cosmic creation and destruction. Thus, every single thing in this way forms a part of the radiant vibration {sphuratta, sphurana) of the light of absolute consciousness. Light and Awareness: The Two Aspects of Consciousness Absolute consciousness understood as the unchanging ontological ground of all appearing is termed 'Prakasa'. As the creative awareness of its own Being, the absolute is called ' Vimarsa '. Prakasa and Vimarsa — the Divine Light of consciousness and the reflective awareness this Light has of its own nature — together constitute the all-embracing fullness (purnata) of consciousness. The Recognition (pratyabhijfia) school of Kashmiri Saivism develops this concept of the absolute which finds its fullest expression in Utpaladeva's Stanzas on the Recognition of God. Even though neither of these two key terms appear in the Stanzas on Vibration or the Aphorisms of Siva, they recur frequently in their commentaries. Thus, although the original formulation of the Doctrine of Vibration differs from the theology of Recognition in this respect, it was extended in the course of its development to accommodate this concept of the absolute as well. This was possible, and quite justified, insofar as the absolute understood in Pratyabhijfia terms does not, as we shall see, differ essentially from that of the Spanda school. We can, as Kashmiri Saivites themselves have done, explain one in terms of the other. The Doctrine of Vibration Prakasa: The Light of Consciousness Prakasa is the pure luminosity' (bhdna) or 'self-showing' that constitutes the essence and ultimate identity (atman) of phenomena. That things appear at all is due to the light of consciousness, and their appearing (avabhasana) is itself this Light which bestows on all things their evident, manifest nature. Established in the light of consciousness everything appears there according to its own specific nature (svabhava). Anything that supposedly does not rest in this Light is as unreal as a sky-flower. 3 Thus, according to Rajanaka Rama, unlike the light of the sun, or any other light, this Light not only makes all things apparent, it is also their ultimate source. 4 Full of its

divine vibration the Light makes all things manifest and withdraws them into itself. This supra- temporal activity characterises it most specifically; devoid of it, it would be no better than an inert physical phenomenon. At the same time, this light is the conjunction (slesa) or oneness (aikdtmya) of its countless manifest forms, 6 and the collective whole (Sampina'ana) of all the categories of existence. The universe is nothing but the shining of the Light within itself. It is the radiant vibration (sphuratta) of this Light, the state (avast hand) in which consciousness becomes manifest. Although the Light shines as all things at all times and hence also makes their diversity manifest, 9 penetrating each object individually as well as collectively, it is not totally 'merged' (magna) or identified with the object so as to suffer any division within itself. Our experience of any object is of the form: T see this': it is not itself an object, but the manifest form the object assumes as a luminous principle of experience. The Light is ever revealed and can never be obscured; objectivity can never cast a shadow on the light of consciousness. The Stanzas on Vibration declare: That in which all this creation is established and from whence it arises is nowhere obstructed because it is unconditioned by [its very] nature. This Light is the highest reality (Paramartha). lity with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**DELAY DUE TO CONFEDERATION DEBIT TYPE MATURATION CREDIT TYPE CONSOLIDATION**

**MATURATION**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \quad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9} + \frac{N_9 dG_9}{dR_9} - \frac{dB_9}{dR_9} > \frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9} + \frac{dM_9}{dR_9}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened**

**eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production Charge Asymmetry, $\Delta Y \equiv Yt - Yt^{-}$, Reflects The Asymmetry In $Tt^{-}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton

particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results

in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely

dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact**

**most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the**

**thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is

nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the

world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease.

Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition

and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the

second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular

chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon

be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining

Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been

Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing

Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

**ABSOLUTE STABILITY**

$$L_9 = \frac{N_9 G_9}{2M_9} \frac{d(9nF_9)}{dR_9} (\geq 0) \quad (9a)$$

$$I_9 = -\frac{N_9 G_9}{2M_9} \frac{d(9nM_9}{dR_9} (\geq 0) \quad (9b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9 > L_9 - I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_9 > L_9 - I_9,$ when $P_9^2/4 > I$ (8a)

$P_9 > L_9 - I_9 + \frac{(I_9 - P_9^2}{(L_9 - I)}$ when $P_9^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to

reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

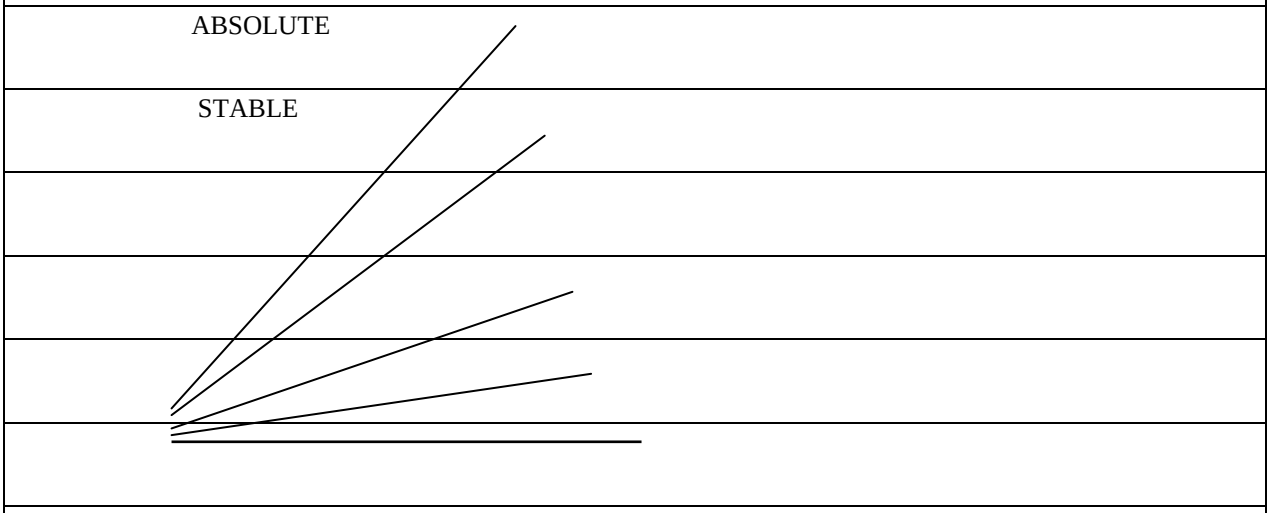



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori

reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

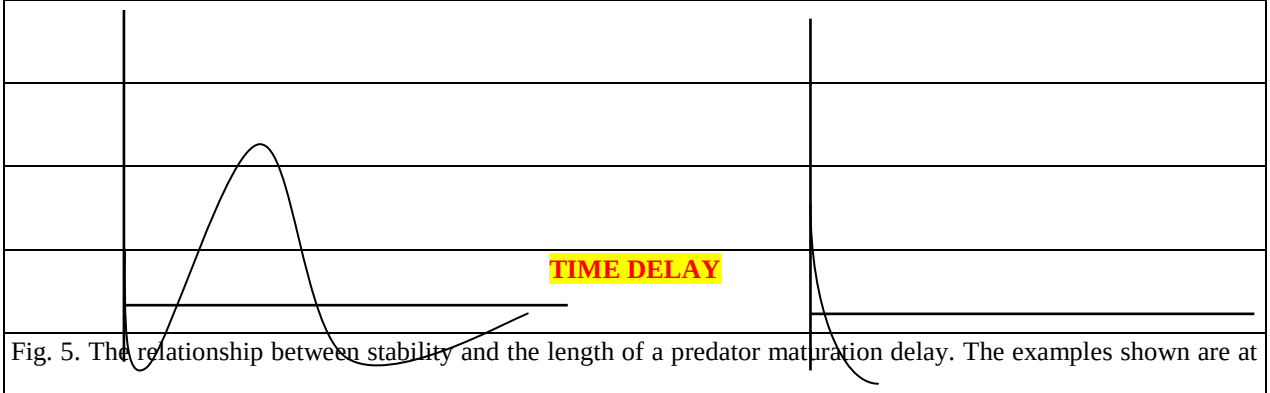



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

**DELAY IN BOTH PREDATOR AND PREY VIS A VIS DEBIT TYPES AND CREDIT TYPES**

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t - T_9)F_9\big(R_9(t - T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t - T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) and the effective prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting

$dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

Systems to which model has been applied also include: **Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI:

http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors.

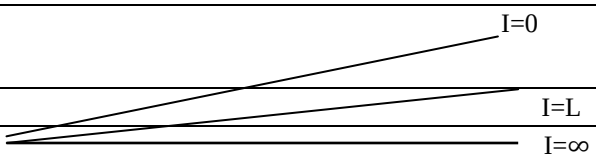
Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, ,**

**Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)**

TOTAL

ABSOLUTELY STABLE

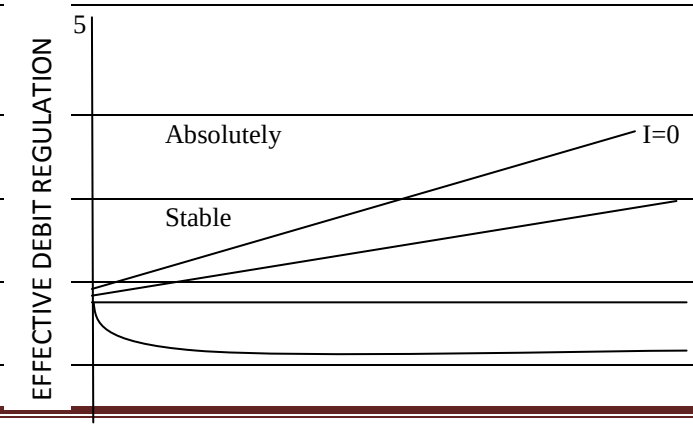



DELAYED DEBIT EFFECT

FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.



-50 10

**DELAYED CREDIT EFFECT (LACK OF BALANCE) CREDIT NON CUMSUMMATION:"**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

The following treatise bears ample testimony and infallible observatory to the fact that how Nunny's model can be applied to various cosmological, general relativity and economic problems.

## ONTOLOGICAL AND EPISTEMOLOGICAL FEATURES OF THE STUDY: OF NUNNY's MODEL

When they are known to be as they are, in the fullest sense of their existence, their phenomenal nature disappears leaving the ground of Being naked and accessible. This approach was validated by a critique of experience. The Vedanta established that space, time and the other primary categories of our daily experience can have no absolute existence. It was therefore necessary to make a distinction between relative truths — that accepted by the precritical common man — and an absolute truth discovered at a higher level of consciousness. The Saiva absolutist 17 rejects any theory that maintains that the universe is less than real. From his point of view a doctrine of two truths, one absolute and the other relative, endangers the very foundation of monism. The Kashmiri Saiva approach is integral: everything is given a place in the economy of the whole. It is equally wrong to say that reality is either one or diverse. Those who do so fail to grasp the true nature of things which is neither as well as both. "We do not" says Abhinavagupta, "base our contention that [reality] is one because of the contradictions inherent in saying that it is dual. It is your approach {paksa) that accepts this [method] While, if [duality and oneness] were in fact [to contradict each other], they would clearly be two [distinct realities]." The Vedantin, who maintains that non-duality is the true nature of the absolute by rejecting duality as only provisionally real, is ultimately landed in a dualism between the real and illusory by the foolishness of his own excessive sophistry (vacafadurvidya). Oneness is better understood as the coextensive unity (ekarasa) of both duality and unity. They are equally expressions of the

absolute. Gopinath Kaviraj says: According to Sarikara, Brahman is truth and Maya is inexplicable (anirvacaniya). Hence the [Advaitin's] endeavour to demonstrate the superiority of Advaita philosophy is turned against his own system. It tarnishes the picture of its philosophical perfection and profundity. He cannot accept Maya to be a reality, therefore his non-dualism is exclusive. The whole system is based on renunciation and elimination and thus is not all-embracing.... By accepting Maya to be Brahman (brahmamayi), eternal (nitya) and real (satyarupa), Brahman and Maya [in the Tantra] become one and coextensive. **The Doctrine of Vibration: An Analysis of the Doctrines and Practices of Kashmir Shaivism Mark S. G. Dyczkowski. Concept of being from Shaivite point of view:** Solutional behaviour and stability analysis of such systems is of paramount importance, notwithstanding the assumptions made which are not farfetched, in that it throws light on the fact whether they could be candidates for yet unsolved problems, providing Authetic ative determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication for the wide amplitudinal ramificatory usage of the thesis propounded. The world is now the other person's world, a foreign world that no longer comes from the self, but from the other. The other person is a "threat to the order and arrangement of your whole world...Your world is suddenly haunted by the Other's values, over which you have no control"(Satre: being and Nothingness) Sartre states that many relationships are created by people's attraction not to another person, but rather how that person makes them feel about themselves by how they look at them. This is a state of emotional alienation whereby a person avoids experiencing **their subjectivity by identifying themselves with "the look" of the other**. The consequence is **conflict**. In order to maintain the person's own being, the person must control the other, but must also control the **freedom of the other "as freedom**". These relationships are a profound manifestation of "bad faith" as for-itself is replaced with the other's freedom. The purpose of either participant is not to exist, but to maintain the other participant's looking at them. This system is often mistakenly called "love", but it is, in fact, nothing more than emotional alienation and denial of freedom. In

essence we are always simulating and selling ourselves at the cost of our freedom. Both as a script writer and participant consciousness, human beings experienceâ complete subjective feelings and sometimes absolute subjective feelings, which if remains unchecked might lead to pathological ramifications. (Italics mine) Perspicuous forbearance, Sophisticated seasoning, Partcipational observation, Background actuality, Existential worldliness, Existential strife, Predicational anteriorities, Character consonance, Ontological consonance, Primordial exactitude, Phenomenological correlates, Accolytish representation, Atrophied asseveration. Anamensial alienisms, Anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orinetatilities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state form the bastion, pillar, post, stylobate and sentinel of the classification scheme and doxa. appropriate and diagnostic, differentiating-discriminating, distinctive-distinguishing, emblematic- especial, essentially exclusive, fixed, idiosyncratic-inborn, inbred, indicative and individual, individualistic-inherent, individualizing and ingrained, inherent, innate, local, marked, native, normal, original, particularly- peculiar, personal, private, proper, regular, representative-singular, special, specific, symbolic, symptomatic- unique these form the classification and stratification measure that is suitable notwithstanding the constancy in many a case, as in gravity, which depends on the mass and distance for two given objects**.** Like Kant and Bergson, Deleuze considers traditional notions of space and time as unifying forms imposed by the subject. He therefore concludes that pure difference is non-spatio-temporal; it is an idea, what Deleuze calls "the virtual". (The coinage refers to Proust's definition of what is constant in both the past and the present: "real without being actual, ideal without being abstract.") Proust, Le Temps Retrouvé, ch. III: see the fourth line from the bottom of this page, or, in English translation, the thirteenth paragraph here: "I began to discover the cause by comparing those varying happy impressions which had the common quality of being felt simultaneously at the actual moment and at a distance in time, because of which common quality the noise of the spoon upon the plate, the unevenness of the paving-stones, the taste of the madeleine, imposed the past

upon the present and made me hesitate as to which time I was existing in. Of a truth, the being within me which sensed this impression, sensed what it had in common in former days and now, sensed its extra-temporal character, a being which only appeared when through the medium of the identity of present and past, it found itself in the only setting in which it could exist and enjoy the essence of things, that is, outside Time. [...] **Nothing but a moment of the past Much more perhaps; something which being common to the past and the present, is more essential than both. [...] a marvellous expedient of nature had caused a sensation to flash to me—sound of a spoon and of a hammer, uneven paving-stones**—simultaneously in the past which permitted my imagination to grasp it and in the present in which the shock to my senses caused by the noise had effected a contact between the dreams of the imagination and that of which they are habitually deprived, namely, the idea of existence—and thanks to that stratagem had permitted that being within me to secure, to isolate and to render static for the duration of a lightning flash that which it can never wholly grasp, a fraction of Time in its pure essence. When, with such a shudder of happiness, I heard the sound common, at once, to the spoon touching the plate, to the hammer striking the wheel, to the unevenness of the paving-stones in the courtyard of the Guermantes' mansion and the Baptistry of St. Mark's, it was because that being within me can only be nourished on the essence of things and finds in them alone its subsistence and its delight. It languishes in the observation by the senses of the present sterilised by the intelligence awaiting a future constructed by the will out of fragments of the past and the present from which it removes still more reality, keeping that only which serves the narrow human aim of utilitarian purposes. But let a sound, a scent already heard and breathed in the past be heard and breathed anew, simultaneously in the present and in the past, real without being actual, ideal without being abstract, then instantly the permanent and characteristic essence hidden in things is freed and our true being which has for long seemed dead but was not so in other ways awakes and revives, thanks to this celestial nourishment "While Deleuze's virtual ideas superficially resemble Plato's forms and Kant's ideas of pure reason, they are not originals or models, nor do they transcend possible experience; instead

they are the conditions of actual experience, the internal difference in itself. "The concept they [the conditions] form is identical to its object."[Desert Islands, p. 36.] A Deleuzean idea or concept of difference is therefore not a wraith-like abstraction of an experienced thing, it is a real system of differential relations that creates actual spaces, times, and sensations. [See "The Method of Dramatization" in Desert Islands, and "Actual and Virtual" in Dialogues. In fact it not as if the style and substance of science and philosophers have been entirely different. look at the conformality and consonance in the Gaussian Theory and Deleuze's statements on identity and the loss of it in the following passage…….That is, identity is a continuous process that finds its limit in exhaustion. As Deleuze will write, "…the logical relation of causality is inseparable from a physical process of signalling, without which it would not be translated into action. By 'signal' we mean a system with orders of disparate size, endowed with elements of dissymmetry; by 'sign' we mean what happens within such a system, what flashes across the intervals when a communication takes place between disparates. The sign is indeed an effect, but an effect with two aspects: in one of these it expresses, qua sign, the productive dissymmetry; in the other it tends to cancel it" (dr, 20). An actualized entity just is such a signal-sign system, remaining in perpetual communication with the virtual multiplicity it actualizes, constituting its own elements and identity. Such is Deleuze's account of what Badiou refers to as the operation of the "count-for-one"… an operation that is immanent to being itself and energetic in character. yet how are we to discern the local nature of ontological situations, the manner in which there is no global, overarching global situation in which all local situations are embedded as if parts of a whole? in part the above passage already responds to this question. if series must be brought into relation to communicate in order for causality to occur, and if the resonance of series is always a matter of chance, of a throw of the dice, like Lacan's objet a where no series enjoys primacy of model over copy, then all we have are local situations without a total situation. Unlike Heidegger's analysis of Dasein where Dasein is always being-in-the-world, Deleuze's simulacra or actualizations are always divergent, enjoying only local relations. Heidegger argues that in order to

understand the being of Dasein we must understand the manner in which Dasein is a being-in-the-world. By contrast, Deleuze's concept of multiplicity provides us with the principle of a new structuralism, of a new local ontology, that allows us to understand the immanent organization of a multiplicity without referring it to an embedding global space. as DeLanda so beautifully puts it in relation to the mathematician gauss, "…when gauss began to tap into these differential resources, a curved two-dimension surface was studied using the old Cartesian method: the surface was embedded in three-dimensional space complete with its own fixed set of axes; then, using those axes, coordinates would be assigned to every point of the surface; finally, the geometric links between points determining the form of the surface would be expressed as algebraic relations between the numbers. but gauss realized that the calculus, focusing as it does on infinitesimal points on the surface itself (that is, operating entirely with local information), allowed the study of the surface without any reference to a global embedding space. basically, gauss developed a method to implant the coordinate axes on the surface itself (that is, a method to implant the coordinate axes on the surface itself (that is, a method of 'coordinatizing' the surface) and, once points had been so translated into numbers, to use differential (not algebraic) equations to characterize their relations. as the mathematician and historian Morris Kline observes, by getting rid of the global embedding space and dealing with the surface through its own local properties, 'gauss advanced the totally new concept that a surface is a space in itself'" (intensive science and virtual philosophy, 11-12).according to the old Cartesian method, we can only outline the properties of a space by relativizing it to a global space in terms of which it is then mapped. by contrast, gauss is able to explore a space in terms of its intrinsic metric and organization as a local space, without referencing it to a whole of which it is conceptualized as a part. **The Image of Thought** "Image of thought "permeates both popular and philosophical discourse. According to this image, thinking naturally gravitates towards truth. Thought is divided easily into categories of **truth and error**. The model for thought comes from the educational institution, in which a master sets a problem and the pupil produces a solution which is either true or false. This image of the subject

supposes that there are different faculties, each of which ideally grasps the particular domain of reality to which it is most suited. In philosophy, this conception results in discourses predicated on the argument that "**Everybody knows..."** the truth of some basic idea. Descartes, for example, appeals to the idea that everyone can at least think and therefore exists. Deleuze points out that philosophy of this type attempts to eliminate all objective presuppositions while maintaining subjective ones. Deleuze maintains, with Artaud, that real thinking is one of the most difficult challenges there is. Thinking requires a confrontation with stupidity, the state of being formlessly human without engaging any real problems. One discovers that the real path to truth is through the production of sense: the creation of a texture for thought that relates it to its object. Sense is the membrane that relates thought to its other. Accordingly, learning is not the memorization of facts but the coordination of thought with a reality. "**As a result, 'learning' always takes place in and through the unconscious, thereby establishing the bond of a profound complicity between nature and mind".** Deleuze's alternate image of thought is based on difference, which creates a dynamism that traverses individual faculties and conceptions. This thought is fundamentally energetic and signifying: if it produces propositions, these are wholly secondary to its development. Philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; even tum tan tum;it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the vorticity of being merges with the positive use of the disjunctive synthesis, and

this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoin der for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for every hum of voices or a single silence for all the silences; a single vacuum for all the vacuousness consciousness should not be said without conjuring if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. The attributable to all states of bodies and states of affairs and the expressible of every proposition So univocity of consciousness means the identity of thé noematic attribute and that which is expressed linguistically and sense-fully. Univovity.We mean that it does not allow consciousness to be subsisting in aquasi state and but expresses in all pervading reality. (Reference: LOS: Deleuze) Deleuze sometimes spoke of getting up behind an author and "creating a monster", as if his reading methodology somehow consisted of a distortion of the philosopher's or author's thought. Certainly there is a case to be made for this with regard to the process of deterritorialization, where something is wrested from a territory and deterritorialized upon a new territory, like the animal paw that is deterritorialized from the earth and reterritorialize on the branch. however, it seems to me that the more interesting aspect of Deleuze's approach to other philosophers and art is not so much his "monstrous creations" (i seldom find them particularly monstrous), but rather their Gaussian or Riemannian style, where he explores them in terms of their own internal organization and metric, without reducing them to something alien such as history, society, biology (reductivism), or the signifier. What we find in Deleuze's approach to phenomena is a Gaussian technique. for instance, take Deleuze's books on cinema, his book on Francis bacon, or his study of sacher-masoch. In the first instance, Deleuze's carefully separates cinema from narrative and the signifier, studying it in terms of its specific organization pertaining to the production of images. in the case of bacon, Deleuze doesn't look for an underlying narrative or "meaning", but instead studies the manner in which bacon composes and

organizes his images and lines, both in terms of their production and actuality, so as to liberate a logic of sensation. finally, in approaching sacher-masoch, Deleuze allows Sacher-Masoch's novels to speak for themselves in terms of their desire and relation to pain, stalwartly refusing to reduce masochism to the complement of sadism. in each case, we have a local exploration of a "space" of multiplicities that is extremely precise. Sean Bowden the priority of events: **Deleuze's logic of sense sean Bowden, the priority of events: Deleuze's logic of sense, Edinburgh university press, 2011, 296pp., $40.00 (hbk), isbn 9780748643646 Reviewed by David Scott, Coppin state university and Gaussian spaces and multiplicities a note on Gaussian spaces and multiplicities posted by larval subjects under uncategorized.** What you "see" is not always what you "get." Many people mistakenly take their own visions literally without "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in the physical sense but in a psychological sense. The archetypes of the collective unconscious are arrayed behind the scenes of current worldwide conditions, of crisis and confusion. They mirror our own states back at us, whether we perceive them as such or interpret them plausibly or not. The noise of ordinary consciousness and beliefs drowns out the signal. Unconsciousness is the background of our ordinary awareness. Our organism is very much at the center of such effects. The organismic source is our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is (=) a meme-Scape. Pre-conceived concepts vie with (e&eb) structures, concepts with images. Like scientists who ignore (e) assumed truths, we leapfrog over (e&eb) our beliefs and personality deficits, claiming (eb) idiosyncratic imagination is (=) literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves (e) us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our

experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images. It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is (=) both. We inherit (e&eb) the structures that make our experience (=) what it is. But the structure itself is "empty," and each human culture "fills" it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about (e&eb) our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from

criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns,

nomenclature, and internal coherence of meaning and representation. n a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar some statements in Vedanta, it is not superfluous or redundant. It is evolutionary. Authetication determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication of the evolutionary process towards understanding is ever present in all, its thematic and discursive form. it is the relentless holding on to a theory this is despicable albeit the theory is not. Look at this paragraph. It is the 'I am the body/mind' belief that gives rise to the 'I am not (e) the world' belief. These two beliefs are (=) co-created."Consciousness projects (e&eb) the appearance of the mind, body and world by taking the shape of thinking, sensing and perceiving." "Attention is Consciousness with an object. When the object vanishes, attention simply remains what it always is, Consciousness.""There is (=) no purpose to (e) meditation. The purpose is already accomplished." "Everything that is experienced is experienced by, through, in and as Consciousness." "The seen cannot be separated from seeing and seeing cannot be separated from (e) Consciousness.""The Reality of any experience is not hidden in the appearance; it is expressed by the appearance." "Once we see that everything is Consciousness... Maya still dances, but it is a dance of love not seduction." There are some truly excellent sections

such as a long one taking, as a starting point, the following quotation from the painter Cézanne: "Everything vanishes, falls apart, doesn't it? Nature is always the same but nothing in her that appears to us lasts. Our art must render the thrill of her performance, along with her elements, the appearance of all her changes. It must give us a taste of her Eternity." Rupert begins: "That statement must be one of the clearest and most profound expressions of the nature and purpose of art in our era." And he goes on for some 14 pages to elaborate on this claim, examining the nature of the 'elements' Existence, appearance and Consciousness and their relationship

Some of the following concepts are modeled; some are not. For those which are not modelled kindly see one of the papers in the series. This note has the axiomatic predication to give a holistic view of the study. Here we talk of the characterstics of systems which satisfy the condition of cosmological constant or any other for that matter or variable or axiomatic predication. . There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit,

facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified

from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. Towards the end of circumvention of typing hundreds of superscripts and subscripts which would be a sardine squash and the concomitant operational difficulties, model is presented in piece meal of nine modules. Logarithms are to be taken in respect of those which incorporate more than one variable in bra-ket. Values of log (ab) and log (a+b) are readily available. Anti logarithm shall be taken at the end to predict, project, prognosticate the value of the variable. This is true for tensors, vectors and other variables too. Affirmational assertion, and explanation justification, statement of vindication, annotational commentary, and explicational glossary for each and every system changes and physical interpretation of results is one thing that is to be with earnest endeavour and feverish and febrile expediency. Editions never ever mean the same and identicality of parameters and this has been explicitly and unmistakably stated in the model apriori itself. Akin and analogous, cognate and concurrent, correspondent and congruent, comparable and complementary, synonymous and duple, tantamount and agnate, commensurate and correlative representation is only to highlight the importance and subterfuge the replication of work. It is my fervent solicitation to kindly bear with me for any lapses, notwithstanding the orchestrated efforts for a paper without any minor errors. Postulation predication, conclusive presumption, differential presuppositions, underscored decidedly axiomatic statement of the statement, equation form the bye word or the watch word in the aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement-argument, articulation, assertion, asseveration, choice of words, commentary and communication, declaration and definition, delivery-diction, elucidation-emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation tone and tenor of stratification. Now, how do find the reaction of systems to these singularities. You do the same thing a boss does for you. "Problematize" the events and see how you behave. I will resort to "pressure tactics". "intimidation of deriding report", or "cut in the increment" to make you undergo trials, travails and tribulations. I am happy to see if you improve your work; but may or may not be sad if you succumb to it and hang yourself! We do the same thing with systems. systems show conducive response, felicitous reciprocation or behave erratically with inner roil, eponymous radicality without and with blitzy conviction say like a solipsist nature of bellicose and blustering particles, or for that matter coruscation, trepidiational motion in fluid flows, or seemingly perfidious incendiaries in gormandizing fellow elementary particles, abnormal ebullitions, surcharges calumniations and unwarranted(you think so but the system idoes not!) unrighteous fulminations. Amount of emotion or affect that is appropriate or expected in an given form of interaction. Particular individuals and diffuse obligations (see c and d) are associated with affectivity, whereas contacts with many individuals (universalistic) in a bureaucracy may be devoid of emotion and characterized by affective neutrality. Affective neutrality may refer to self discipline and the deferment of gratification (eg Weber's spirit of capitalism). In contrast, affectivity may be associated with expressing emotions. Adams and Sydie also refer to affective neutrality being associated with ego control (p. 15). (Parsons: Wikipedia). So the point that is made here is "like we problematize the "events" to understand the human behaviour we have to "problematize" the events of systems to understand their behaviour. This statement is made in connection to the fact that there shall be creation or destruction of particles or complete obliteration of the system (blackhole evaporation) or obfuscation of results. Some systems are like "inside traders" they will not put signature atoll! How do you find they did it! Anyway, there are possibilities of a CIA finding out as they recently did! So we can do the same thing witn systems to. This is accentuation, corroboration, fortification, .fomentatory note to explain the various coefficients we have used in the model as also the dissipations called door. In the bank example we have clarified that various systems are individually conservative, and their conservativeness extends holisticallytoo.that one law is universal does not mean

there is complete adjudication of nonexistence of totality or global or holistic figure. Total always exists and "individual" systems always exists, if we donot bring Kant in to picture! For the time being let us not! Equations would become more ensorcelled and frenzied..... philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". In consciousness research, two rival sets of theories can be recognized: (A) Scientific material interpretations of consciousness are based on axioms that view consciousness in the context of highly advanced intentional processing of information in which subject-object relations evolve, and (B) humanistic interpretations of consciousness are based on axioms that view consciousness in the context of, say, "centered pulsations" that enable a conscious agent to act from his or her center of awareness. In this paper I will argue for the selection of axioms that favor humanistic interpretations of consciousness. The Pursuit of Autonomy Interdisciplinary Observations to Human Consciousness: Wautischer, Helmut (1993) the Pursuit of Autonomy. Interdisciplinary Observations to Human Consciousness: Social Neuroscience Bulletin, 6 (4). pp. 52-56.Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality; **CONSCIOUSNESS AND ITS UNIQUENESS June 12, 2012 at 8:40pm (See Deleuze Logic of sense, Wikipedia and Stanford encyclopedia for more details)** Atrocious contrivance and device, gratificational primogeniture, calamitous dodge and expedient, depraved, destructive, disastrous, execrable gambit and gimmick, great solace and succor, iniquitous, injurious, loathsome, low, maleficent, malevolent machination and maneuver, promethaleon of candor, frankness, honesty, honor, ingenuousness, innocence, openness, reality, sincerity, truthfulness, spiteful, stinking, ugly, unpleasant, unpropitious play and ploy, progenitor of jurisprudence and circumspection, racket and ruse, savvy, scam, stratagem, subterfuge, tactic and wile, forthrightness, honesty, truthfulness embodied and personified is the structural predisposition and dispensation of the complex. Nobody need have to take the Brahman-Anti Brahman agency to be a recalcitrant, repugnant, and refractory proposition. God gives and also takes away. That is the point made. ‘

And of course lesser gods, and God doth follow them and they follow God. That is why they are lesser gods. And wield the power on lesser mortals like us. When we mean gratification and deprivation we mean celestial Brahman Anti Brahman and Terrestial Brahman and Anti Brahman agencies producing such happiness or sadness by any means such as (blending- amalgamation, adulterant-adulteration, amalgam-amalgamation, blend- combination,

composite- compound, debasement-denaturant, fusion-hybrid, intermixture-reduction, amalgam- mixture, composite- compound, fusion-mishmash, aggregate- alloy, amalgamation-blend, commixture-composite, composition-compost, conglomerate-fusion, goulash-medley, mishmash-stew, blend-coalition, commixture-compound, federation-heating, integration-junction, melting-merger, synthesis-unification, assortment- combination, adulteration-alloy, assimilation-association, batter-blend, brewed combine, composite compound, concoction confection, conglomeration cross of money, orientation(dis), penury creation, murder, mayhem plunder, pillage, apocalypse, Armageddon in various forms and permutation and combination of Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness: with each action accounted and showing balance in the nature's general ledger as also actionlessness:e). All these are recorded in Nature's general ledger. Neuron DNA encapsulates all the actions passions, actions, interactions, transactions or the lack of it at various levels of individual, collective and cosmic levels. In celestial Brahman Anti Brahman ledger there is cessation of dualities. One more point to be made is if 5=3+2, then we can say 5 gormandizes 3 or 5 gobbles up 2. Quintessentially such a statement arises out of the difference of LHS and RHS term wise. It must be remembered despite all theoretical abstractions and generalizations be it proton mass or Higgs Boson, or topology of a space, it is notwithstanding the fundamentality of nature is just a number. Just everything in Physics and mathematics is just a number. Let us assume I buy rice in Bangalore (District) worth 5k.g. I can always divide 5k.g.s of rice in categories of Bangalore

(South), Bangalore (west), Bangalore (East) and Bangalore (South). On similar basis, there are billions of neutrons, protons and concomitant interactions which would satisfy the given numerical value. Example of gravity above makes things clear. Everything atleast is a function of time and someplace or some other coordinates, subordinates or superordinates. Mass of proton is 2.47 does not mean total mass of all protons is 2.47. It is this argument that is used in the classification scheme. By the very nature, self the unmonitoring but the witness consciousness agency identifies itself with something, be it a profession, an actor on the screen, when it does not have a identity. Such examples are legion, a Charlie Chaplin in "The Kid" or Raj kapoor in "Jagte Raho". These are people who just wanted to live and living without being put under surveillance every moment is the greatest part of their life. This is "Absolute Subjectivity" of Shiva. Such a state is always craving for identification for in pure consciousness all identities are lost. It is this "identitylessness" or having very minimal identity is what pervades most in India and this is due to expansion of individual consciousness. When that happens, be it education, or wealth, the "identity' becomes so well entrenched, they become the exact opposite of "pure consciousness"," impure consciousness" (Shakti). This is a state of "Absolute or relativistic objectivity". There was this case of a lady who went to such an extent of planning, she was sure that whatever she does must work out. When it did not, she lapsed in to "identity crisis". Such thing as this might happen in well developed countries where the competition is extremal and to be achieved at any cost. Realisation that one is not the doer but the witness consciousness is just like watching a movie in which one has acted. It is in this sense "advaita", "non duality' seems to be of cardinal and preeminent importance for the analytical treat ment of "identity' and "identity lessness", both happen in space time. "Multiverses" in this sense is also like a different film of an actor released at different theaters and it is in this sense we state that that every object in space time is a multiverse. You must have totally immersed in doing a "string theory" problem or watching a "film", oblivious of surroundings. Essentially creation (space time) also is what is happening in mind and honed in to you billion times. Identification of anything with a designation or signification is one that leads to

association within spacetime, which would only have pathological consequences and detrimental ramification as stated above or positive implications of "happiness" because of the identity itself. Search or 'show" has to go on! It gives a sense of consummation. a watching pot never boils; donot go to check nature; you are disturbing it; it will never show its true colours bhakti and Virakthi, gratification and deprivation, raga and dwesha is essential to live in space. It is the quantum interest thereof that produces energy to live .gratification and deprivation which form the bastion, pillar, post and stylobate and sentinel of all human actions depends on tams, rajas, and sattva of individual or insentient object. Verily tamas, rajas and satva are the characterstics of space itself, nay the behavioural pattern and attitudinal orientation. There are peaceful places like Vienna or Sringeri and there are violent places like others. Space is warped and woofed by the Akshara, the quantum information of primordial syllable Om. cosmic general ledger is the one vibration of which is brahman; quantum information possessions thereof form the sentinel of all transactions (Gargi to Yajnavalkya: my interpretation)adhocism, nonseqientiality of events, anti sequentiality of events, simultaneity of phenomenon, concurrent happenings , improbability of phenomenon happening itself are characterstics of brahman-anti brahman (Terrestial and celestial) are the cause of non predictability of happenings like in black hole, or in the areas of operation of brahman-anti brahman (Terrestial and celestial). brahman AntiBrahman actions, interactions, decisions, circumspection, adhocism, satisfies sensitive to initial conditions; topologically mixing; dense periodic orbits, the axiomatic predications and postulation alcovishness of chaos theory. Like gravity is a characteristic of space, so is tamas, rajas, sattva, with spaces knotted together by quantum information (for models see below) .accidents are preplanned actions of Terrestial and celestial Brahman AntiBrahman. All actions are perpetrated by them, be it train on fire or death. Belief systems play an important role in the determination of tamas, rajas, and satva state of human beings. There is neither cause nor effect. All actions are performed by Terrestial and celestial Brahman anti Brahman and it is off the record. psychic energy is negative equivalent to –mc^2, so that second law of thermodynamics is satisfied on a holistic basis vis-à-vis material energy

ParaBrahaman is infinite individual, collective and cosmic consciousnesses are infinite (concomitant general ledgers) for they are in the evolution mode with respect to each other in that order. Later emanates from the former as transactions emanate from the conservative state of assets and liabilities. Is not zero true for infinite conservative systems inside space and time? zero is the only consummation, consolidation and conservation. if fifty out of seventy members come out of the room at whatever time you come out, talk about the grand limbo of oblivion and hibernation of lifeless world you are going to and if you are not a celebrity, then it is not coincidence, it is coordination. Nature is playing the same trick on you. That a generalised law is there and you can measure does not make things less good albeit recalcitrant, refractory and rogue particles occur time and again, by and large. Entering in to all living beings through his constant part on the new moon night, he is born therefrom next morning. Absolute ParaBrahaman (zero) projects himself again and again on the screen of individual consciousness which appears of depairing despondent world. Some where the balancing act is done on karma. cosmic general ledger reflects upon itself to keep all transactional ties conservative (Brihadarnkayaka Upanishad: 1.5 14) for models see next paper through its divine power ( i interpret man creates god; it is attribution and ascription of various Shaktis would the formation of god takes place like space and time) the self assumes (eb) the form of a deity man can contemplate and venerate, even though siva, the pure subject, can never in fact be an object of meditation (adhyeya). Until we realise our true identity with Sankara, he is worshipped and conceived to be a reality alien to ourselves. While we are in the realm of creation, he too is a creation or mode or appearing of the absolute, manifest to us in meditation, through his freedom as an eternal, omniscient being. There is no gulf between the created and the uncreated creator: nothing in reality, although an object of knowledge, ceases to be Siva: this is the reason why meditation [on this or that aspect] of reality bestows its fruit. The world of the senses and mind appears to the well awakened (suprabuddha) as a theophany, an eternal revealing of god in his creation. The doctrine of vibration declares that "there is no state in word, meaning or thought, either at the beginning, middle or end, which is not Siva." 59 to utter any word is, in

reality, to intone a sacred formula.every act is a part of Siva's eternal cosmic liturgy, every movement of the body a ritual gesture (mudras), and every thought, god's thought. By what path are you not attainable? What words do not speak of you? in which meditation are you not an object of contemplation? What indeed are you not, o lord? 61 spandashakti, which accounts for the appearing of all things, is also the means by which deity in its many varied forms appears to man. Ksemaraja concludes: the ultimate object of worship of any theistic school differs not from the Spanda principle. The diversity of meditation is due solely to the absolute freedom of Spanda. 62 Sankara is not only the supreme object of devotion; as the static polarity of the absolute, he is the inner reality which holds together its Siva and Sakti outer manifestations. Phenomena are patterns of cognitions projected onto the surface of self-luminous siva-consciousness. There they become apparent, directly revealed to consciousness according to their manifest form. Siva is accordingly symbolized as the ground or surface of awareness, smooth and even like a screen (samabhittitalopama). 64 inscribed on this screen (kutfya) are the countless manifest forms which appear within it rendering it as diverse and beautiful as a fossil ammonite (falagrama). 65 Siva is the sacred ground upon which the cosmic mandalas are drawn, the absolute surface of inscription which bears the mark (cihna) of the universe. Abhinava writes: the variety of this world can only be manifest if the highest lord, who is essentially the pure light of consciousness, exists; just as a surface is necessary for a picture. if external objects were perceived in isolation then, because 'blue' and 'yellow', etc., are self-confined and the perceptions [we have of them] refer to their objects alone and so are insentient, mute and dumb in relation to one another . . . how would it be possible to be aware that an object is variegated? but just as depths and elevations can be represented by lines on a smooth wall, and we perceive [a female figure and think], similarly it is possible to be aware of differences in the variegated (contents of experience) only if all the diverse perceptions are connected together on the one wall of the universal light of consciousness. 66 Siva is the perfect artist who, without need of canvas or brush, paints the world pictures. The instant he imagines it, it appears spontaneously, perfect in every respect. The colours he uses are the varying shades

and gradations of his own Spanda energy and the medium his own consciousness. Intractable disaffect, Mutinous obdurate, Rebellious recalcitrant, Seditious refractory, Insubordinate insurgent, antagonistically antipodean, discrepant incompatible, Contumacious coadjutant, Conflicting converse, Unyielding obdurate, the antagonistic and anachronistic dispositional ties of Anti Brahman are in fact the sine qua non of Brahman only (See Vishwaroopa of Mahavishnu, which contains the polarised positives on RHS and negatives on LHS). The universe is coloured with the dye of its own nature (svabhava) by the power of Siva's consciousness (a//). Rajanaka rama says: homage to him who paints the picture of the three worlds, thereby displaying in full evidence his amazing genius {pratibha)\ to Sambhu who is beautiful with the hundreds of appearances laid out by the brush of his own unique, subtle and pure energy. 68 analogously, at the microcosmic level, all the cognitions and emotions, etc., which make up the individual personality form the outward flow of essentially introverted consciousness. They are specific pulsations (vise$aspanda) or aspects of the universal pulsation (samanyaspanda) of pure t consciousness. at the lower level, within the domain of maya, they represent the play in the fettered soul of the three primary qualities (gunas) or feeling-tones' which permeate to varying degrees his daily experience. These are: 1) sattva — the quality of goodness and luminosity which accompanies blissful experience both aesthetic and spiritual. 2) Rajas — the passion or agitation which oscillates between the extremes of 'light' and 'darkness' and characterises inherently painful experiences. 3) Tamas — the torpor and delusion which accompany states of inertia and ignorance. 69 the liberated soul recognises that these three are the natural and uncreated powers of pure consciousness. for him they are manifest respectively as: 1) Sankara 's power of knowledge (jnana) — the light of consciousness (prakasa); 2) the power of action (kriya) — the reflective awareness of consciousness (vimarsa)\ 3) the power of maya — which does not mean here the world of diversity, but the initial subtle distinction which appears between subject and object in pure consciousness. 70(doctrine of vibration: mark S.G. Dyczkowski) .to "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new

concepts and meanings that could not be included in the old system. at the heart of western metaphysics, for example, derrida finds the opposition between "speech" and "writing." this binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior in derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. The postmodern turn in philosophy: theoretical provocations and normative deficits by Steven best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html. You cannot see beyond consciousness for that is the limit of your storage of information or the basis for understanding further information. Turiya is empty bucket; it is from which prajna the measure of  arises; turiya measures the number of quantum information stores in other three states of waking, dreams, and dreamless deep sleep (Manduka Upanishad; my interpretation for models next paper).let there be path where angels tread: it is a force never born never dies the brahman an anti brahman verily there are parallel universes in (eb) in the ambit of space and time; the society of sinners and the sinned with the grammar of evil; parallel universes contain(e) multiverses in them; all physics, philosophy,  break down(e) in these universes; astronomical gratification ,deprivation, glorification, mortification, projections, axes, dimensions, rotations, foldings, concurrences, motions, shocks, are the axiomatic predications of multiverses and parallel universes; death is when self (the witness consciousness; register of individual general ledger) leaves the body; ostensibly parallel seem to have (e) an objective; parallel universes donot (e) have objectives; their paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common

attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, affection- appreciation, awareness- discernment, emotion-feeling, gut reaction, heart insight, intuition, judgment, keenness, perceptiveness, rationale, sensation, sense, sensitiveness, sensitivity, sentiment, susceptibility, taste, vibes, admiration, aesthetic sense, affection, appraisal, assessment, attraction, awareness, cognizance, commendation, comprehension, enjoyment, esteem, estimation, grasp, high regard, knowledge, liking, love, perceptional realization, recognitional regard, relish, respect, responsiveness, sensibility, sensitiveness, sensitivity, sympathy, understanding, valuation, acknowledgment, airing, baring, , confessional defenselessness, denudation- denunciation, disclosure, display, divulgence-exhibition, hazard, introduction, jeopardy, laying open, liability, manifestation, nakedness, openness, peril, presentation, publicity, revelation, risk, showing, susceptibility , susceptiveness,susceptivity, unfolding, unmasking, unveiling, vulnerability, vulnerableness, action, affection, appreciation and ardor, behavior, capacity, compassion, concern, cultivation and culture, delicacy, discernment and discrimination, emotion, empathy, faculty, fervor, fondness, heat, imagination, impression and intelligence, intensity and intuition, judgment, keenness- palpability, passion, pathos, pity, reaction, refinement, sensibility, sensitivity, sentiment, sentimentality , sharpness, spirit, sympathy, tangibility, taste, tenderness, understanding, warmth, fond memories, hearts and flowers, homesickness, longing, pining, reminiscence, ,schmaltz, sentimentality wistfulness, yearning satiation is only ostensible; their job is to get the job done; their business is to mind other's business ;all laws of physics and finance breakdown; it is antilaw,piscatorial Piratish and bubonic bucaneerishness; there is free transactions between in multiverses; when reflected on the screen of consciousness they appear as extant existential universe. Transactions are astronomical and gratification

deprivation, glorification and mortification incalculable; they have no barriers of region, religion, language which they create for others. theirs objectives are: corruption criminality, debasement debauchery, degradation evil immoral lewdness licentiousness, perversion sinfulness, vice viciousness, vileness wickedness, improbity corruption, atrocity decadence, degeneration degradation, depravity evil, immorality impurity, infamy iniquity, looseness lubricity, perversion profligacy, sinfulness turpitude,vice viciousness, vulgarity wickedness,disgrace bad reputation, abasement abuse, baseness black eye, blemish blur, brand comedown,contempt contumely, corruption culpability, debasement debasing, defamation degradation, derision disbarment, discredit disesteem, disfavor dishonor, disrepute disrespect, humbling humiliationignominy ill repute, infamy ingloriousness, meanness obloquy, odium opprobrium, pollution ,put-down reproach, scandal scorn, slander slight, slur spot, stain stigma, taint tarnish, turpitude venality, enormity Horribleness , abomination atrociousness, atrocity crime and the positive aspectionalities thereof .let there not be any incursions where evil mercenaries in the garb of mendicants thrive with adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatious gainsaying or conformingly pugnacious Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, execrable and goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness . statement for

Mephistophelean mercenaries philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality. Negative energy stabilising wormholes are considered candidates for (er) darkmatter. An analysis of negative energy might be a useful ingredient of the thoughts of travel at speeds more than the velocity of light. Peirce explains that an "imperfect continuum" possesses "topical singularities, or places of lower dimensionality" which break continuity, while "the parts of a perfect continuum have the same dimensionality

as the whole." (CP 4.642) Circular continua without singularities are time and space. (CP 6.210-212) By using math as his lens, Peirce envisioned singularities or lower dimensional holes in the imperfect cosmologic continuum, while maintaining that a probabilistic law of chance exists that entropy decreases **(Wikipedia).** On the side of philosophy**,** quite explicitly, Deleuze provides a critique of Husserl's "sleight-of-hand" change from a static categorical phenomenology to one supposedly more dynamically constitutive of reality. This a clear cut **chagrined circumvention**, **contraventional curbing**, **defeated disgruntlement**, **dissatisfactional downer**, **dragged failure**, **fizzled foil**, **grievance hindrance**, **impedimental irritation**, letdown, nonfulfillment, nonsuccess, obstruction, old one-two, **resent mental setback**, unfulfillment of the earlier theories makes Sartre contend that human existence is a conundrum whereby each of us exists, for as long as we live, within an overall condition of nothingness (nothing-ness)—that ultimately allows for free consciousness. But simultaneously, within our being (in the physical world), we are constrained to make continuous, conscious choices. It is this dichotomy that causes anguish, because choice (subjectivity) represents a limit on freedom within an otherwise unbridled range of thoughts. Subsequently, humans seek to flee our anguish through action-oriented constructs such as escapes, visualizations, or visions (such as dreams) designed to lead us toward some meaningful end, such as necessity, destiny, determinism (God), etc. Thus, in living our lives, we often become unconscious actors—Bourgeois, Feminist, Worker, Party Member, Frenchman, Canadian or American—each doing as we must to fulfill our chosen characters' destinies. However, Sartre contends our conscious choices (leading to often unconscious actions) run counter to our intellectual freedom. Yet we are bound to the conditioned and physical world—in which some form of action is always required. This leads to failed dreams of completion, as Sartre described them, because inevitably we are unable to bridge the void between the purity and spontaneity of thought and all-too constraining action; between the being and the nothingness that inherently coincide in our self. Deleuze insists on the absolute non-resemblance between what conditions (metastable, pre-individual plane of singularities) and what is conditioned

(actualized or the individuated 'thing'). Only then do "the conditions of the true genesis become apparent," Deleuze writes (Logic of Sense 105). Without addressing this issue, Deleuze's borrowings from Simondon, Lautman and structuralism remain untethered. Because of Bowden's decision to completely sidestep Deleuze's overt negotiations with phenomenology, his otherwise sophisticated commentary suggests a pedagogical narrowness. If, the omission of phenomenology strengthens Bowden's focus on the ontological priority of the event over substantial ontologies via the genealogy of figures and movements Deleuze outlines, it also incompletely articulates the problem of the sense-event as Deleuze formulates it. We have lost any idea of whom Deleuze takes to be his primary adversary, against whom he stakes out his position. This worry is most clearly manifested in Bowden's description of Deleuze's ontological prioritizing of events as offering a "transcendental ontology." For Bowden Deleuze's transcendental ontology defines events that "are ontologically prior to substances, 'all the way down'" (82). But Bowden provides a second definition: it is that the priority of events over worldly individuals means that the world of events "is not something external to the conditions of knowledge" (69). **SEAN BOWDEN the Priority of Events: Deleuze's Logic of Sense Sean Bowden, the Priority of Events: Deleuze's Logic of Sense, Edinburgh University Press, 2011, 296pp., $40.00 (hbk), ISBN 9780748643646. Reviewed by David Scott, Coppin State University Weltschmerz, agony, apprehension, blues, depression, dread, mid-life crisis, misgiving, nervousness, uneasiness, angst were expressed with such effulgent words…… aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement, argument, articulational assertion, asseverational choice of words, commentary and communication, declaration and definition, delivery, diction, elucidation, emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation about the human predicament is brought out in Being and Nothingness. Being and Nothingness offers a critique of Sigmund Freud's theory of the unconscious, based on the claim that consciousness is essentially self-conscious. Sartre also argues that**

**Freud's theory of repression is internally flawed. According to Sartre, in his clinical work, Freud encountered patients who seemed to embody a particular kind of paradox—they appeared to both know and not know the same thing. In response, Freud postulated the existence of the unconscious, which contains the "truth" of the traumas underlying the patients' behavior. This "truth" is actively repressed, which is made evident by the patients' resistance to its revelation during analysis. Yet what does the resisting if the patients are unaware of what they are repressing? Sartre finds the answer in what Freud calls the "censor". "The only level on which we can locate the refusal of the subject," Sartre writes, "is that of the censor.** The Holographic Principle, as conceived in current physics, applies to fields, and perhaps, even to more elementary entities, called strings and branes. These more elementary entities remain hypothetical at the time of this writing. There are many layers between the level of fields and that of neurons. We know that neurons are surrounded by a surface phospholipid membrane, which supports an electrochemical process that is fundamentally necessary to human experience. As the information supported by the membrane surface according to the Holographic Principle is proportional to its area, we should expect to find that the neural correlates of consciousness most clearly associated with those parts of the neuron that have the have a high proportion of neuronal surface area, and, in particular, receptive surface area. Dendrites account for an average of about 90% of the neuron is receptive surface area (Wong, 2002), and so would be expected to be the most important neural structures in information processes. We find this to be true **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science (Wikipedia).** Discordant-discrepant, dissonant-hard-line, implacable, incompatible, incongruous- inconsistent, inexorable-inflexible, inharmonious-intransigent, opposed, reluctant, unappeasable, uncompromising, unfriendly some of these expatiations and enumerations might sound, a more generalised interpretations is given by a Vedantin. Look at this arresting passage: The teacher continued to be silent. When addressed a second and third time he said: "I am teaching, but you do not follow. The Self is silence." The undetermined and unthinkable character of the Brahman is

a consequence of the absolute's eternal and immutable nature. To concede the existence of a real universe is, from the Vedantin's point of view, to posit the existence of a reality apart from the Brahman. Nor can we simply identify a real universe with the absolute unless we are prepared to compromise its unchanging, absolute status. The criterion of authenticity is immutability. Reality never changes; only that which is less than real can appear to do so. Reality is constant in the midst of change. What this means essentially is that there is change although nothing changes. This impossible situation is reflected in the ultimate impossibility of change itself. That which does not exist prior to its changing and at the end, after it has changed, must be equally non-existent between these two moments. Although the world of change appears to be real, it cannot be so. Change, according to the Vedantin, presupposes a loss of identity. Reality cannot suffer transformation; if it were to do so, it would become something else and the real would be deprived of its reality. The immortal can never become mortal, nor can the mortal become immortal. The ultimate nature of anything cannot change. Change of any sort is merely apparent (vivarta); the world of change and becoming is a false super- imposition (adhyaropa, adhyasa) on the absolute. In cosmic terms, the mistake (bhranti) consists of the supposition that the real Brahman is the unreal universe and the unreal universe is the real Brahman. In microcosmic terms, it is the mistake of falsely conceiving the body, mind or even one's personality to be the Self. In the same way as the image of a snake is falsely superimposed on a rope, similarly the universe is falsely projected onto the real substratum, the Brahman. Ignorance is not merely a personal lack of knowledge, but a cosmic principle. As such it is called "Maya," the indefinable factor (anirvacaniya) that brings this mistake in identity about. The reality status of this cosmic illusion is also indefinable: on the one hand it is not Brahman, the sole reality; on the other hand it is not absolutely non-existent like a hare's horn or the son of a barren woman. Brahman is the source of world appearances only in the sense of being their unconditioned ground or essential nature. The universe is false not because it has no nature of its own but because it does have one. Just as the illusion of a snake disappears when one sees that it is nothing but a rope, similarly cancellation (badha) of the empirically real occurs

when the absolute reality of the Brahman is realised. Thus, according to Vedanta, appearance implies the real, while the real need not imply appearance. To appear is essentially to appear in place of the real, but to be real is not necessarily to appear. All things -exist because the absolute exists. It is their Being. Thus the very existence of phenomena implies their non-existence as independent realities.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^\infty \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9 \left( U_9 - V_9 \int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9 \right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9 \int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$then given the constraints of $f(T_9), (A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 -$$

$$V_9 M_9 -]e^{-T_9\lambda_9} = 0$$

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9 Z_9 \sin(Z_9 T_9) - (A_9 M_9 + Y_9) + (A_9 M_9 + X_9)\cos(Z_9 T_9) = 0$$

$$(A_9 M_9) Z_9 - M_9 Z_9 \cos(Z_9 T_9) + (A_9 M_9 - X_9)\sin(Z_9 T_9) = 0$$

Quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High

Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) **Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP** This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) **Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP** In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference: Journal of Physics: Conference Series 70 **Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows:** System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol

Where $A_9 = U_9 - V_9$Using (A2.3), substitute for $\boldsymbol{Z}_9^2$ and $Z_9$in (A2.2):

$W_9^2 (X_9^2 - 3 M_9 A_9 X_9 - M_9^2 (X_9 - Y_9) + 2 M_9^3 A_9 + 2 A_9^2 M_9^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_9$=1-cos ($Z_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (X_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ X_9 - Y_9 /2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \leq W_9 \leq 2$ *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (X_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\ M_9)^2\ (2A_9 M_9 - (X_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9 M_9^2 - 3M_9 X_9\right) + W_9\left(3M_9 A_9^2 - 2A_9 X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9+Y_9\right)$$

$$=2M_9^2W_9^2(M_9+\mathrm{A}_9)+M_9A_9^2W_9+\left(2A_9M_9-X_9+Y_9\right)(A_9W_9+M_9W_9^2$$

$$+(2-W_9)((2M_9X_9+M_9Y_9)W_9+A_9(X_9+Y_9)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range 0 $W_9$ 2.Defining

$Q_9\left(W_9,A_9\right)=a_9W_9^2+c_9$ *where* $a_9b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

To Heidegger writings we shall compose the model. Following is not still written model able form. For one of the models please see any one of the papers under the appellation: Ïn Excited Reveries

In his later writings Heidegger drops the term 'phenomenon' in preference for the verbal form 'phainesthai' in order to emphasize even more the actuality or presentational property of Being. Explaining this new form of the term he writes: "Being disclosed itself to the ancient Greeks as 'physis'. The etymological roots 'phy-' and 'pha-' designate the same thing: 'phyein', the rising-up or upsurge which resides within itself as 'phainesthai', lighting-up, self-showing, coming-out, appearing-forth." Heidegger contrasted his notion of phenomenon with semblance (Schein) and with appearing (Erscheinung). In the case of semblance a thing can show itself as that which it is not, as when fool's gold shows itself to be gold. The ancients always allied semblance with non-being. Heidegger points out, however, that semblances are grounded in

showings, and so does Abhinava. Both Heidegger and Abhinava consequently maintain that all semblances have a real basis and are to be treated as instances of phenomena along with the so-called real showing or manifestation of non-deceptive objects. So Heidegger states that: 'how- ever much seeming, just that much being'. Thus self-showing or appearing defines Being as phenomenon, but this definition of Being is as yet incomplete. Being is not only self-showing but 'logos' which Heidegger explains means 'discourse' (Rede) in the sense of 'apophansis': 'letting-be-seen'. Phenomenology, which according to Heidegger is the only correct study of Being, means 'letting-be-seen-that-which- shows-itself. This is true of Saiva Paramadvaita as well. The reality of the world demands recognition; we are forced to accept the direct presentation of the fact of our daily experience. As Abhinava says: "if practical life, which is useful to all persons at all times, places and conditions were not real, then there would be nothing left which could be said to be real" A thousand proofs could not make 'blue' other than the colour blue. The reality of whatever appears in consciousness cannot be denied. Objects appear; they do not cease to do so by a mere emphatic denial. The manifestation of an entity in its own specific form is a fact at one level of consciousness; it is real. The appearing of the same entity in the same form but recognised to be a direct representation of the absolute is also a fact, but at another level of consciousness. It is no more or less real than the first. 'As is the state of consciousness, so is the experience,' says Abhinava. Although the nature of the absolute is discovered at a higher level of consciousness, The Doctrine of Vibration nonetheless it presents itself to us directly in the specific form in which we perceive things; otherwise there would be no way in which we could penetrate from the level of appearing to that of its source and basis. Abhinava writes: Real is the entity iyastu) that appears in the moment of direct perception (sak$atkara), that is to say, within our experience of it. Once its own specific form has been clearly determined one should, with effort, induce it to penetrate into its pure conscious nature. All things are known to be just as they present themselves. The concrete actuality of

being known (pramiti), irrespective of content, is itself the vibrant (spanda) actuality of the absolute. Liberating knowledge is gained not by going beyond appearances but by attending closely to them. "The secret," Mahesvarananda says, "is that liberation while alive (jivanmukti) is the profound contemplation of Maya's nature." No ontological distinction can be drawn between the absolute and its manifestations because both are an appearing (dbhasa), the latter of diversity and the former of 'the true light of consciousness which is beyond Maya and is the category Siva'. Those who have attained the category of Pure Knowledge above Maya and have thus gone beyond the category of Maya, see the entire universe as the light of consciousness . . . Just as the markings [on a feather] are nothing apart from the feather, the feather [is nothing apart from] them, similarly, when the light of consciousness is manifest, the whole group of phenomena is manifest as the light of consciousness itself. Within the sphere of Maya, every entity's 'own nature' (svabhdva) corresponds to its specific manifest form. Accordingly it is defined as that which distinguishes it from all else and from which it never deviates. Above the sphere of Maya, that is, above the level of objectivity, is the domain of the subject. At this level, everything is realised to be part of the fullness of the experience r and hence no longer bound by the conditions which impinge on the object. Here the part is discovered to be the whole, that is, consciousness in toto. In this sphere beyond relative distinctions, the yogi realises that (all) the categories of existence are present in every single category. m The yogi experiences every individual particular as the sum total of everything else. He recognises that all things have one nature and that every particular is all things. This is the 'essence' (sard) or co-extensive unity (samarasya) of all things. Integral Monism of Kashmiri SaivismWe have established that reality is manifest according to how [and the degree in which] the freedom of consciousness reveals it and that [this freedom] is the womb of all forms. Just as 'sweetness' is present in its entirety in every atom of the sugarcane, so each and every atom [of the universe] bears within itself the emanation of all things. This is

the level of consciousness in which the absolute reflects on itself realising to its eternal delight and astonishment (camatkara) its own integral nature. The reality of the world of diversity is not denied, but experienced in a new mode of awareness free of time and space in the eternal omnipresence of the Here and Now. [Phenomenal forms of awareness] such as 'this [exists]', born of the colouring [imparted to the absolute] by the limitations engendered by the diversifying power of time {kalakalana) also emanate within the Supreme Principle. There [at that level], Fullness {purriata) is the one nature [of all things] and so everything is omnipresent; otherwise, associated with division (khancjlana), the Fullness [of the absolute] would not be full. The content of absolute consciousness consists of diverse appearings (abhdsa) which, because they are manifest through it in this way, do not compromise the wholeness of consciousness. Everything we perceive is a momentary collocation of a number of such manifestations which combine together like 4 a row of altar lamps' (dipavalT) to form the single radiant picture of the universe. The individual objects which constitute the universe are specific collocations of such 'atomic' appearings. Together they form a single unified particular which appears according to its own defining features (svalaksand). A jar, for example, consists of a number of appearances such as 'round', 4 fat\ 'earthen', 'red', etc., which together discharge a single function (arthakriya), in this case, that of carrying the appearance 'water'. They unite with each other much as the scattered rays of a lamp come together when focused, or as the various currents of the sea together give rise to waves. Atomic appearings can combine in any number of ways, provided that they are not contrary to one another as established by the dictates of natural law (niyati). An appearance of 'form', for example, cannot combine with that of 'air'. Insofar as they share a common basis (samanyadhikaranya), a given cluster of appearances appears as a single whole. This common basis is the most prominent member of the group; the appearance 'jar' is such in the example quoted above. Any one appearance in a cluster may assume a more important or subordinate role.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_9 \square - b_9 \square 0$$

And

$$b_9^2 \square 4a_9 c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9\, in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\ X_9$ ,

$Let\ Y_9 = \ y_9\ X_9$And try the substitution $A_9 = X_9/(2 + A_9\ y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + \ n_9 y_9)$ we have

$$W_9^2\, y_9 \left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^2 \left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9\,X_9\left[\frac{X_9\left(1+n_9y_9\right)}{M_9^2\left(2+n_9y_9\right)^2}+2\,y_9\left(n_9-1\right)+\frac{X_9^2\left(9+y_9\right)}{M_9^2\left(2+n_9y_9\right)}\right]=0.\ (A2.12)$$

Since $n_9$ and $A_9$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9=\left(2+n_9y_9\right),$ *the condition*$(A2.99?)$*can be written as*

$$\left[\frac{X_9\left(E_9-9\right)}{E_9^2}\right]^2\left[\frac{X_9^2}{M_9^2}-4\,y_9\,E_9^2\frac{X_9}{M_9^2}+4y_9\left(n_9-2\right)E_9^3\right]\ \text{[illegible]}0.$$

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9/W_9^2)$ can be expressed as

$$2\,E_9\left[y_9E_9+\left(y_9\,E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}\right]>\frac{X_9}{M_9^2}>2E_9\left[y_9E_9-y_9E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}]$$

dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique

for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) **Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP** Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) **Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP** Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a

cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) **Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP** Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service

placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) **Dynamic Service Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP** Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by $H*(X,(\wedge*TX)\otimes g)$, where X is the target hyperk\"ahler manifold. We further give a mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) **A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van**

**Raamsdonk1 and Nathan Seiberg2** study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. **Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP** Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy **The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP** Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the

present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. **ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP** Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition. Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb) the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537

© 2016 American Physical Society **Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP** Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant ($\Delta\alpha/\alpha$) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and $\Delta\alpha/\alpha$ to (e&eb) this conclusion. However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts $z\lesssim 3$, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the $\Lambda$ cold dark matter model value ($j\Lambda CDM=1$); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in $\alpha(z)$, linked to (e&eb) a linear coupling ($\zeta$) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order $|\zeta|\sim 10-4$, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in $\Delta\alpha/\alpha$ (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society **The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP** Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation

rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void

Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems, authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society **Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP** Authors measure the star formation properties of (e) two large samples of

galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the **star formation** rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a **function of environment**. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]**.** Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size $\sqrt{n} \times \sqrt{n}$? The problem is that $\sqrt{n}$ time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time $O(\sqrt{n})$ for $d \geq 3$, or $O(\sqrt{n}\log 5/2n)$ for $d = 2$. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search

tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ∗This work was mostly done while the author was a **PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship and by CIAR**. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 **Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP** Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can

lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 **Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP** Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. **Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP** Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e)

sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) **Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP** The relationship between the foundations of mathematics and (e&eb) physics is a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb) flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling

field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) **Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP** The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar

and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph] (or arXiv: 1503.05600v1 [math-ph] for this version) **Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP** This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs

mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) **Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP** Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and

derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to **the fine structure constant**. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) **Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP** Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor goes to 0 as the (e&eb) time goes to 0, the

big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph] (or arXiv: 1306.4613v1 [quant-ph] for this version) **Effects of gauge theory based number scaling on geometry Paul Benioff STOP** In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x,

outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph]

(or arXiv: 1211.3381v2 [quant-ph] for this version) **Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP** Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value,

\int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances, the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) **Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP** The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include

(e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) **Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP** Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) **New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP** This work is based on (e) the field of reference frames based on (e) quantum

representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects:Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc) Journal reference:Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) **A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul B**

**ABSOLUTE STABILITY**

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 ? $W_9$ ? 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9-9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2-2y_9\frac{(E_9+n_9^2y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right)-4y_9(E_9-n_9)$$ (A2.14)

And has the solution

$$\frac{X_9}{M_9^2}>\frac{E_9}{(E_9-9)}[y_9(E_9+n_9^2y_9)+[y_9^2(E_9+n_9^2y_9)^2+4y_9(E_9-9)(E_9-n_9)]^{9/2}]$$

Given $E_9>n_9$this lower limit to $X_9/M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9/M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9<n_9$ $(ie.,2A_9M_9^2<X_9-Y_9)$then (A2.9)and (A2.8)show that $Q_9$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2}>2E_9[y_9E_9-(y_9E_9(9+y_9)(E_9-n_9)]^{9/2}]$$

And

$$E_9>n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2/2M_9(X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$ and $X_9$ and $Y_9$ by $W_9^2$ the dependence upon *M* disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

**Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability

of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e) long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007 STOP. David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer P** examine the possibility of coherent reversible information transfer between solid-state superconducting qubits and (e&eb) ensembles of ultracold atoms. Strong coupling between these systems is mediated by (e&eb) a microwave transmission line resonator that interacts near resonantly with (e&eb) the atoms via (e&eb) their optically excited Rydberg states. Solid-state qubits can then be used to (e) implement rapid quantum logic gates, while (e) collective metastable states of (e) the atoms can be employed for (e) long-term storage and (e&eb) optical readout of quantum information. **Reversible state transfer between superconducting qubits and (e&eb) atomic ensembles David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer Phys. Rev. A 79, 040304(R) – Published 14 April 2009 STOP.** Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures have been proposed. An

attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. **Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier** demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in (eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. **Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP.** In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown

to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , **Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally** frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. **Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg** propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states

into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. **Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011 http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan** proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 **Qubit transfer between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller** propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. **Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys.**

**Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich** report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. **Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi** investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. **Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev** investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer.

Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. **L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP.** In recent years **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P have** investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P** apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys. Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. **Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S.**

**Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP.** Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert **Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe** realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.**220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2** [quant-ph] for this version) **Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP.** Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. **Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker** describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue

from over (e&eb) a thousand plant and animal samples. **NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP.** Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this behavior, enhancing and accelerating (eb, eb+) the growth of correlations. **Michael Kastner** gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. **Michael Kastner** tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) **Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP.** A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. **Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E. Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91,**

**075435 – Published 26 February 2015 STOP.** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still

problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to

(e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP** Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that

(eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035

**Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime**

**C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013**

**STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and

cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP** Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex

Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) **A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen**

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9 and$ separate (A3.1)into real and imaginary parts

$$Z_9^2 + \left(M_9 + V_9\right) Z_9 \sin\left(Z_9 T_9\right) - \left(U_9 M_9 + Y_9\right) + (M_9\left(U_9 + V_9\right) - X_9)$$

$$\mathrm{X}\cos\left(Z_9 T_9\right) - 2M_9 V_9 cos^2\left(Z_9 T_9\right) + M_9 V_9 = 0$$

$$(M_9 + U_9)\, Z_9 - (M_9 + V_9)\, Z_9\, cos\left(Z_9 T_9\right) + \left(M_9\left(U_9 + V_9\right) - X\right)$$

$$\mathrm{X}\sin\left(\left(Z_9 T_9\right) - M_9 V_9 \sin\left(Z_9 T_9\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 =$ $1 - cos\, Z_9\, T_9$and $A_9 = U_9 - V_9$we obtain

$$4\, M_9^2 V_9^2 W_9^4 + 2\, M_9 V_9\, W_9^3[(M_9 - V_9)^2 + A_9\left(M_9 + V_9\right) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2\left[M_9 - V_9\right)^2\left(2\, A_9\, M_9 + Y_9 - X_9\right) + \left(2A_9\, M_9 - X_9\right)^2 - A_9\left(M_9 - V_9\right)\left(2A_9\, M_9 - X_9\right)$$

$$+ M_9 V_9\,(A_9^2 + 4\left(X_9 + Y_9\right))]$$

$$+ W_9\left[A_9^2 + 2X_9\right)\left(2\, A_9\, M_9 - X_9\right) + 2A_9 Y_9\,(M_9 + V_9) - M_9 A_9^3] + A_9^2\left(X_9\, M_9 - Y_9\right) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9\left(2, U_9 V_9\right) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting cos $(Z_9 T_9) = -1\ in\ (A3.3)$

We have applied the model to Kashmiri Shaivism. Whose elaborate rituals is enough to make a indivual raise from individual consciousness to cosmic consciousness which we have called Cosmic Consciousness.

It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) **Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP** The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic

Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** The group of unimodular transformations on the roots of polynomials over GF(2) is considered, and those polynomials with (e&eb) symmetries in (eb) the unimodular group are identified. The cross-correlation function between (e&eb) two maximum-length linear shift register sequences of (e) the same degree is shown to be (=) computable as an explicit linear transformation, in (eb) matrix form, on (e&eb) either one of the sequences, regarded as (=) a vector. The underlying vector space is (=) the "cyclotomic algebra," generated by (e) the cyclotomic cosets, or by what Gauss termed the **"periods" of the cyclotomic equation**. A variety of numerical examples are worked in detail. **Information Sciences Volume 1, Issue 1, December 1968, Pages 87-109 Theory of transformation groups of polynomials over GF(2) with applications to linear shift register sequences Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more link. Solomon W. Golomb** stop Authors study the P−V criticality and phase transition in (eb) the extended phase space of charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of (e) the system and its conjugate quantity is (=) the thermodynamic volume of the black hole. The black holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb) no P−V criticality and phase transition appear, while (e) for the black holes with a spherical horizon, even when the charge of the black hole is absent, t(eb) he P−V criticality and the small black hole/large black hole phase transition will appear, but it happens only in d=5 dimensions; when the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to have (e) the P−V criticality

and (e&eb)the small black hole/large black hole phase transition, there exists (ebb) an upper bound for the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the **Gauss-Bonnet coefficient** and Q is (=) the charge of the black hole. Authors calculate the critical exponents at the critical point and find that (eb) for all cases, (eb) they are the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09(2013)005 Cite as: arXiv:1306.6233 [gr-qc] (or arXiv:1306.6233v4 [gr-qc] for this version) **P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in IR but with (e&eb) improved UV behaviors. Here, authors study an IR modification which breaks ``softly'' (e) the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Authors obtain the black hole and cosmological solutions for **``arbitrary'' cosmological constant** that represent (eb) the analogs of the standard Schwarzschild-(A)dS solutions which can be asymptotically (A)dS as well as flat and authors discuss their thermodynamical properties. They also obtain solutions for (e) FRW metric with (e&eb) an arbitrary cosmological constant. Authors study its implication to (e&eb) the dark energy and find that it seems to be consistent with (e&eb) current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP.** The Einstein equations with (e&eb) a negative cosmological constant admit (e) black hole solutions which are asymptotic to (e) anti-de Sitter space. Like black holes in asymptotically flat space, these solutions have (e) thermodynamic properties including (e) a characteristic temperature and intrinsic entropy equal to (=) one quarter of the area of the event horizon in Planck units. There are however some important differences from (e) the asymptotically flat case. A black hole in anti-de Sitter space has (e) a minimum temperature which occurs when (e) its size is of the order of the characteristic radius of (e) the anti-de Sitter space. For larger black holes the red-shifted

temperature measured at (eb) infinity is greater. This means that such black holes have (e) positive specific heat and can be in stable equilibrium with (e&eb) thermal radiation at a fixed temperature. It also implies (eb) that the canonical ensemble exists for asymptotically anti-de Sitter space, unlike the case for asymptotically flat space. One can also consider the microcanonical ensemble. One can avoid (e) the problem that arises in asymptotically flat space of (e) having to put the system in (eb) a box with unphysical perfectly reflecting walls because (e) the gravitational potential of anti-de Sitter space acts as a box of finite volume. Communicated by A. Jaffe **Communications in Mathematical Physics December 1983, Volume 87, Issue 4, pp 577-588 Thermodynamics of black holes in anti-de Sitter space S. W. Hawking, Don N. Page STOP** In previous work authors introduced the concept of black-hole entropy, which they identified with (e&eb) the surface area of the black hole in question expressed in units of (e) the Planck length squared. Authors suggested that the appropriate generalization of the second law for a region containing a black hole is (=) that the black-hole entropy plus (+) the common entropy in the black-hole exterior never decreases. Here authors establish the validity of this law for the infall of an entropy-bearing system into (e&eb) a much larger and more massive generic stationary black hole. To do this authors determine (eb) a general lower bound for the increase in black-hole entropy, and an upper bound for the entropy of the system, while allowing for quantum effects at each stage. In passing authors show that the generalized second law is a statistical law which becomes (=) over-whelmingly probable in the limit of a macroscopic system. Authors also consider briefly more general situations. Finally, we give two simple examples of predictions made by the generalized second law for (e) black-hole formation processes. Received 17 September 1973 DOI:http://dx.doi.org/10.1103/PhysRevD.9.3292 **Generalized second law of thermodynamics in black-hole physics Jacob D. Bekenstein Phys. Rev. D 9, 3292 – Published 15 June 1974 STOP** The Lovelock gravity consists of (e) the dimensionally extended Euler densities. The geometry and horizon structure of (e) black hole solutions could be quite complicated in (eb) this gravity; however, authors find that some thermodynamic quantities

of the black holes like the mass, Hawking temperature and entropy have (e) simple forms expressed in terms of (e) horizon radius. The case with black hole horizon being a Ricci flat hypersurface is (=) particularly simple. In that case the black holes are (=) always thermodynamically stable with (e&eb) a positive heat capacity and their entropy still obeys (e&eb) the area formula, which is no longer valid for (e) black holes with positive or negative constant curvature horizon hypersurface. In addition, for black holes in (eb) the gravity theory of Ricci scalar plus a 2n-dimensional Euler density with (e&eb) a positive coefficient, thermodynamically stable small black holes always exist in (eb) D=2n+1 dimensions, which are absent in (eb) the case without (e) the Euler density term, while (e) the thermodynamic properties of the black hole solutions with (e&eb) the Euler density term are qualitatively similar to those of black holes without (e) the Euler density term as D>2n+1. Copyright © 2004 Published by Elsevier B.V. **Physics Letters B Volume 582, Issues 3–4, 4 March 2004, Pages 237–242 A note on thermodynamics of black holes in Lovelock gravity Rong-Gen CAI STOP** There are a number of similarities between black-hole physics and (e, eb, =) thermodynamics. Most striking is the similarity in the behaviors of black-hole area and (e&eb) of entropy: Both quantities tend to increase (eb+) irreversibly. In this paper we make this similarity the basis of (e) a thermodynamic approach to black-hole physics. After a brief review of the elements of the theory of information, authors discuss black-hole physics from the point of view of information theory. Authors show that it is natural to introduce (eb) the concept of black-hole entropy as (=) the measure of information about (e&eb) a black-hole interior which is (=) inaccessible to an exterior observer. Considerations of simplicity and consistency, and dimensional arguments indicate (eb) that the black-hole entropy is equal to (=) the ratio of the black-hole area to the square of the Planck length times a dimensionless constant of (e) order unity. A different approach making use of the specific properties of Kerr black holes and of concepts from information theory leads to (eb) the same conclusion, and suggests (eb) a definite value for the constant. The physical content of (e) the concept of black-hole entropy derives from (e) the following generalized version of the second law: When common entropy goes down a

black hole, (eb) the common entropy in (eb) the black-hole exterior plus the black-hole entropy never decreases. The validity of this version of the second law is supported by (e) an argument from information theory as well as by several examples. Received 2 November 1972 DOI:http://dx.doi.org/10.1103/PhysRevD.7.2333 **Black Holes and Entropy Jacob D. Bekenstein Phys. Rev. D 7, 2333 – Published 15 April 1973 STOP** The effects of power-law Maxwell field on the Van der Waals like phase transition of higher-dimensional dilaton black holes are probed in detail. It is shown that the Smarr relation gains corrections due to (e) the effects of (e) both the power-law Maxwell field and (e&eb) the dilaton field while (e) thermodynamic volume is exactly the same as (=) that of Einstein-Maxwell-dilaton black holes. Authors successfully derive the analytic solutions of (e) critical point and carry out (e&eb) some check to ensure that these critical quantities are (=) positive. It is shown that the constraint on the parameters turns out to be (=) 0<\alpha^2<1, which is tighter than that in (eb) the non-extended phase space. It is also shown (eb) that these critical quantities and (e&eb) the ratio P_cv_c/T_c are affected by (e&eb) the power-law Maxwell field. Moreover, critical exponents are found to coincide with those of other AdS black holes, showing (eb) the powerful influence of mean field theory. Comments: references added Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1601.05500 [gr-qc] (or arXiv:1601.05500v2 [gr-qc] for this version **Effects of power-law Maxwell field on the Van der Waals like phase transition of higher dimensional dilaton black holes Jie-Xiong Mo, Gu-Qiang Li, Xiao-Bao Xu STOP** A consistent and unified picture for (e) critical phenomena of charged AdS black holes in f(R) gravity is drawn in this paper. Firstly, we investigate the phase transition in canonical ensemble. Authors derive the explicit solutions corresponding to (e&eb) the divergence of CQ. The two solutions merge into (e&eb) one when (e) the condition Qc=−13R0−−−√ is satisfied. The curve of specific heat for Q<Qc has two divergent points and can be divided into (e&eb) three regions. Both the large radius region and the small radius region are (=) thermodynamically stable with (e&eb) positive specific heat while (e) the medium radius region is unstable with (e&eb) negative specific heat. However, when Q>Qc, the specific heat is

always positive, implying (eb) the black holes are locally stable and no phase transition will take place. Secondly, both the T−r+ curve and T−S curve f(R) AdS black holes are investigated and they exhibit Van der Vaals like behavior as the P−v curve in the former research. Critical physical quantities are obtained and they are consistent with (=) those derived from the specific heat analysis. Authors carry out numerical check of Maxwell equal area law for the cases Q=0.2Qc, 0.4Qc, 0.6Qc, 0.8Qc. The relative errors are (=) amazingly small and can be negligible. So the Maxwell equal area law holds for (e) T−S curve of f(R) black holes. Thirdly, authors establish geometrothermodynamics for (e) f(R) AdS black hole to examine the phase structure. It is shown that the Legendre invariant scalar curvature R would diverge exactly where (e) the specific heat diverges. To summarize, the above three perspectives are consistent with (e&eb) each other, thus providing (eb) a unified picture which deepens the understanding of (e) critical phenomena of (e) f(R) AdS black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1602.01251 [gr-qc] (or arXiv:1602.01251v1 [gr-qc] for this version) **A consistent and unified picture for critical phenomena of f(R) AdS black holes Jie-Xiong Mo, Gu-Qiang Li, Yu-Cheng Wu STOP** The state equation of the charged AdS black hole is reviewed in (e&eb) the T−r plane. Thinking of the phase transition, the T−S, P−V, P−v graphs are plotted and then the equal area law is used (e) in the three cases to get (eb) the phase transition point (P, T). The analytical phase transition point relations for (e) P-T of charged AdS black hole has been obtained successfully. By comparing the three results, authors find that the equal area law possibly cannot (e) be used directly for P−v plane. According to the T−S, P−V results, they plot (e&eb) the P−T−Q graph and find that for a highly charged black hole a very low temperature condition is required for (e) the phase transition. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI:10.1140/epjc/s10052-015-3641-0 Cite as: arXiv: 1503.07658 [gr-qc] (or arXiv:1503.07658v1 [gr-qc] for this version) **A note on Maxwell's equal area law for black hole phase transitionShan-Quan Lan, Jie-Xiong Mo, Wen-Biao Liu STOP** Recently, extended phase space thermodynamics of Lovelock AdS black holes has been of great interest. To provide

insight from (e) a different perspective and gain (e) a unified phase transition picture, non-extended phase space thermodynamics of (e) (n+1)-dimensional charged topological Lovelock AdS black holes is (=) investigated detailedly in (eb) the grand canonical ensemble. Specifically, the specific heat at (eb) constant electric potential is calculated and phase transition in the grand canonical ensemble is discussed. To probe the impact of the various parameters, they utilize (e) the control variate method and solve (e) the phase transition condition equation numerically for the case k=1, −1. There are two critical points for (e) the case n=6, k=1 while there is only one for other cases. For k=0, there exists no phase transition point. To figure out the nature of phase transition in the grand canonical ensemble, authors carry out an analytic check of the analog form of (e) Ehrenfest equations proposed by Banerjee et al. It is shown that Lovelock AdS black holes in (eb) the grand canonical ensemble undergo (e&eb) a second order phase transition. To examine the phase structure in (eb) the grand canonical ensemble, authors utilize the thermodynamic geometry method and calculate both the Weinhold metric and Ruppeiner metric. It is shown that for both analytic and graphical results that the divergence structure of (e) the Ruppeiner scalar curvature coincides with (=) that of the specific heat. Research provides (eb) one more example that Ruppeiner metric serves as (=) a wonderful tool to (e) probe the phase structures of (e) black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C75 (2015) 5, 211 DOI:10.1140/epjc/s10052-015-3434-5 Cite as: arXiv: 1503.01956 [gr-qc] (or arXiv: 1503.01956v1 [gr-qc] for this version **Non-extended phase space thermodynamics of Lovelock AdS black holes in grand canonical ensemble Jie-Xiong Mo, Wen-Biao Liu STOP** To provide an analytic verification of the nature of phase transition at the critical point of P−V criticality, the original expressions of Ehrenfest equations have been introduced directly. By treating the cosmological constant and (e&eb) its conjugate quantity as (=) thermodynamic pressure and volume respectively, authors carry out analytical check of (e) classical Ehrenfest equations. To show that approach is universal, authors investigate not only higher-dimensional charged AdS black holes (e&eb), but also rotating AdS black holes. Not only are the examples

of Einstein gravity shown, but also the example of modified gravity is presented for (e) Gauss-Bonnet AdS black holes. The specific heat at constant pressure CP, the volume expansion coefficient α and the isothermal compressibility coefficient κT are found to diverge exactly at the critical point. It has been verified that both Ehrenfest equations hold at (eb) the critical point of P−V criticality in (eb) the extended phase spaces of (e) AdS black holes. So the nature of the critical point of P−V criticality of AdS black holes has been demonstrated analytically to be (=) a second-order phase transition. These results are consistent with (e&eb) the nature of liquid-gas phase transition at the critical point. In this sense, extant research would deepen the understanding of (e) the relations of AdS black holes and (e&eb) liquid-gas systems. Moreover, successful approaches to introduce (eb) the original expressions of Erhenfest equations directly into (e&eb) black hole phase transition research demonstrate again (eb) that black hole thermodynamics is closely related (e&eb) to classical thermodynamics, which allows (eb) us to borrow techniques from (e) classical thermodynamics to investigate the thermodynamics of black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 89, 084057 (2014) DOI: 10.1103/PhysRevD.89.084057 Cite as:arXiv: 1404.3872 [gr-qc] (or arXiv:1404.3872v1 [gr-qc] for this version) **Ehrenfest scheme for P−V criticality of higher dimensional charged black holes, rotating black holes and Gauss-Bonnet AdS black holes Jie-Xiong Mo, Wen-Biao Liu STOP** Aiming at a unified phase transition picture of (e) the charged topological black hole in (eb) Ho\v{r}ava-Lifshitz gravity, authors investigate this issue not only in (eb) canonical ensemble with (e&eb) the fixed charge case but also in grand-canonical ensemble with (e&eb) the fixed potential case. Authors firstly perform the standard analysis of the specific heat,(e&eb) the free energy and the Gibbs potential, and (e&eb) then study its geometrothermodynamics. It is shown that the local phase transition points (eb) not only witness the divergence of (e) the specific heat, but also witness the minimum temperature and the maximum free energy or Gibbs potential. They also witness the divergence of (e) the corresponding thermodynamic scalar curvature. No matter which ensemble is chosen, the metric constructed can

successfully produce (eb) the behavior of the thermodynamic interaction and (e&eb) phase transition structure while other metrics failed to (e) predict the phase transition point of (e) the charged topological black hole in former literature. In grand-canonical ensemble, authors have discovered the phase transition which has not been reported before. It is similar to (e) the canonical ensemble in which the phase transition only takes place when k=−1. But it also has its unique characteristics that the location of the phase transition point depends on (eb) the value of potential, which is different from (e) the canonical ensemble where (e) the phase transition point is (=) independent of the parameters. After an analytical check of Ehrenfest scheme, authors find that the new phase transition is (=) a second order one. It is also found that the thermodynamics of (e) the black hole in Horava-Lifshitz gravity is quite different from that in (eb) Einstein gravity. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP1310(2013)056 DOI:10.1007/JHEP10(2013)056 Cite as: arXiv:1404.2497 [gr-qc] (or arXiv:1404.2497v1 [gr-qc] for this version) **A unified phase transition picture of the charged topological black hole in Horava-Lifshitz gravity Jie-Xiong Mo, Xiao-Xiong Zeng, Gu-Qiang Li, Xin Jiang, Wen-Biao Liu STOP** To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case $\beta\to\infty$ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness condition of (e) entropy. It is shown

that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure authors have discovered. It is shown that Legendre invariant thermodynamic scalar

curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws, leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions, geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged AdS black holes by considering the case of the RN-AdS black holes

surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase transition. For the case $\omega_q=-2/3$ authors derive analytic expressions of (e) critical physical quantities, while (e) for cases $\omega_q\neq-2/3$ appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this

version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=) chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy

Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed

quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions, while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma?**

All this consists of the act of consciousness alone", says Vasubandhu, "because unreal entities appear, just as a man with defective vision sees unreal hair or a moon, etc."He points to dreams as examples of purely subjective constructs which appear to be objective realities. The apparent reality dreams possess is not derived from any concrete, objective world, but merely from the idea of objectivity. While the Yogacara does not say that an idea has, for example, spatial attributes, it does have a form manifesting them. While he agrees with the Saiva idealist that

appearances have no independent existence apart from their appearing to consciousness, he maintains that for this reason they are unreal. The creativity of consciousness consists in its diversification in many modes having apparent externality; it is not a creation of objects. While the Kashmiri Saivite agrees that the world is pure conscious- ness alone, he maintains that it is such because it is a real creation of consciousness. The effect is essentially identical with the cause and shares in its reality. Matter and the entire universe are absolutely real, as 'congealed' (sty ana) or 'contracted' (samkucita) forms of consciousness. "This God of consciousness", writes Ksemaraja, "generates the universe and its form is a condensation of His own essence (rasa)" m By boiling sugarcane juice it condenses to form treacle, brown sugar and candy which retains its sweetness. Similarly, consciousness abides unchanged even though it assumes the concrete material form of the five gross elements. The same reality thus abides equally in gross and subtle forms. Consequently no object is totally insentient. Even stones bear a trace (vasana) of consciousness, although it is not clearly apparent because it is not associated with the vital breath (prana) and other components of a psycho-physical organism. Somananda goes so far as to affirm that physical objects, far from being insentient, can only exist insofar as they are aware of themselves as existing. The jar performs its function because it knows itself to be its agent. Indeed, all things are pervaded by consciousness and at one with it and hence share in its omniscience. Thus, Siva, Who perceives Himself in the form of physical objects, is the one ultimate reality. "The jar knows because it is of my nature", writes Somananda, "and I know it because I am of the jar's nature. I know because I am of Sadasiva's nature and He knows because He is of my nature; Yajnadatta [knows] because he is of Siva's nature and Siva [knows] because He is of Yajnadatta's nature". Integral Monism of Kashmiri Saivism Everything in this sense is directly perceived by absolute conscious- ness, and this direct perception (pratyaksa) unifies the knowable into a single, undivided whole. This is the central concept behind a doctrine originally expounded by Narasimha called 'the non-dualism of direct perception' (pratyaksadvaita). This states that consciousness is essentially perceptive and that its perception of all things operates throughout the universe. Insofar

as phenomena are clearly evident (sphufa) to us, everything is directly perceived by absolute consciousness, with which our individual consciousness is identical. This direct perception unfolds everywhere; the one true reality, it is alone and without companion or rival (nihsapatna). Even though it remains one, it can, by its very nature, perceive distinctions (bheda) between one entity and another, without this engendering any division within it. We distinguish between two entities in empirical terms on the basis of their mutual exclusion (anyonyabhava). The relative distinction {bheda) between them is essentially the perceived difference between their respective characteristics. Despite this difference they are united within the purview of a single cognition insofar as they are equally both manifest appearances. This cognition is the undivided essence (rasa) or 'own nature' (svabhava) of both. Encompassed by the 'fire of consciousness', there is no essential difference between them. Just as when an emerald and ruby reflect each other's light, the ruby is reddish-green and the emerald greenish-red, similarly everything is connected with everything else as part of the single variegated (vicitra) cognition of absolute consciousness. Mahesvarananda writes: The Supreme Lord's unique state of emotivity (asadharanabhava) is the outpouring of pure Being (mahasatta). It is manifest as the brilliance (sphuratta) of the universe which, if we ponder deeply, [is realized to be] the single flavour (ekarasa) of the essence of Beauty which is the vibration of the bliss of one's own nature. In this way all things are in reality one although divided from the one another sharing as they do the 'single flavour' (ekarasa) of the pure vibration of consciousness. Kashmiri Saiva Realism Kashmiri Saivism as a whole has been variously called a form of 'realistic idealism', 'monistic idealism', ‘idealistic monism' and 'concrete monism'. It is easy to understand why Kashmiri Saivism is The Doctrine of Vibration said to be 'idealistic' and 'monistic', but in what sense is it also 'realistic*? The answer to this question is of no small importance in trying to under- stand the central idea behind its metaphysics and the fundamental importance of the concept of Spanda, in this seemingly impossible marriage between monistic idealism and pluralistic realism. The Kashmiri Saiva approach understands the world to be a symbol of the absolute, that is, as the manner in which it presents itself to us. Again we can

contrast this view with that of the Advaita Vedanta. The Advaita Vedanta understands the world to be an expression of the absolute insofar as it exists by virtue of the absolute's Being. Being is understood to be the real unity which underlies empirically manifest separateness and as such is never empirically manifest. It is only transcendentally actual as 'being-in-itself. The Kashmiri Saiva position represents, in a sense, a reversal of this point of view. The nature of the absolute, and also that of Being, is onceived as an eternal becoming (satatodita), a dynamic flux or Spanda, 'the agency of the act of being'. It is dentified with the concrete actuality of the fact of appearing, not passive unmanifest Being. Appearance (abhasa) alone is real.ppearing (prakasamanatva) is equivalent to the fact of being (astitva). Ksemaraja writes in his commentary on the Stanzas on Vibration: Indeed, all things are manifest because they are nothing but manifestation. The point being that nothing is manifest apart from manifestation The absolutely unmanifest, from this point of view, can have as little existence as the space in a lattice window of a sky-palace. Nay, even less, because even that space can appear as an imagined mage manifest within consciousness. Everything is real according to the manner in which it appears. Even an illusion is in this sense real, insofar as it appears and is known in the manner in which it appears. The empirical and the real are identical categories of thought. As Abhinava says: Thus this is the supreme doctrine (upanisatf), namely that, when- ever and in whatever form [an entity] appears, that then is its particular nature. Perhaps at this stage a brief comparison with Heidegger's ideas might prove to be enlightening and not altogether out of place. According to Heidegger's phenomenology of Being, reality is intelligible in a two-fold Integral Monism of Kashmiri Saivism manner as 'phenomenon' and 'logos'. Heidegger defines what he means by 'phenomenon' as: "that-which-shows-itself. The manifest . . . phenomena are then the collection of that which lies open in broad daylight or can be brought to the light of day — what the Greeks at times implicitly identified as 'ta onta' (the things-which-are)".

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb) heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna**

**Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version) **Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb) strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel

period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar (=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI: 10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills (SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear

differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black**

**holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-

relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb)

the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this

version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstate thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11 August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of

entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.** Characterizing the behaviour of strongly coupled

quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at λ=λc, the correlation length diverges as ξ∼|λ−λc|−ν, while (e) the gap between the ground state and first excited state vanishes as (e) ΔE∼ξ−z∼|λ−λc|zν. Exactly at the critical point λ=λc, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as ω∼|k−k0|z. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-**

**equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004 STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element.

**Model has been applied for the following results:**

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| **Searches for Sterile Neutrinos with the IceCube Detector M. G. Aartsen et al. (IceCube Collaboration) Phys.**<br><br>**Rev. Lett 117, 071801 – Published 8 August 2016 Physics logo See Viewpoint: Hunting the Sterile Neutrino** |
| **(1) The IceCube neutrino telescope at the South Pole has measured the atmospheric muon neutrino spectrum as (=) a function of zenith angle and (e&eb) energy in the approximate 320 GeV to 20 TeV range, to search for (e) the oscillation signatures of (e) light sterile neutrinos.** |
| **(2) No evidence for anomalous νμ or ν¯μ disappearance is observed in (eb) either of two independently developed analyses, each using one year of (e) atmospheric neutrino data.** |
| **(3) New exclusion limits are placed on (e&eb) the parameter space of (e) the 3+1 model, in which muon antineutrinos experience (e) a strong Mikheyev-Smirnov-Wolfenstein-resonant oscillation.** |
| **(4) The exclusion limits extend to (e) sin22024≤0.02 at Δm2∼0.3 eV2 at the 90% confidence level.** |
| **(5) The allowed region from (e) global analysis of appearance experiments, including (e) LSND and MiniBooNE, is excluded at approximately the 99% confidence level for (e) the global best-fit value of \|Ue4\|2. Received 5 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.117.071801 © 2016 American Physical Society** |
| **Synopsis: Searching for Majorana Neutrinos August 16, 2016The KamLAND-Zen collaboration has run the most sensitive search to date for a radioactive decay that could reveal whether neutrinos are Majorana fermions. Synopsis figureKamLAND-Zen Collaboration** |
| **(6) Are neutrinos and (=) antineutrinos the same particles, as hypothesized by (e) physicist Ettore Majorana?** |
| **(7) If true, this "Majorana" nature of neutrinos could explain (eb) why neutrinos have (e) mass and why the Universe has (e) more matter than antimatter.** |
| **(8) Using an underground neutrino detector in Japan, the KamLAND-Zen collaboration has carried out the most sensitive search to date for a process called neutrinoless double-beta decay, which could confirm (eb) Majorana's hypothesis.** |
| **(9) The results set stringent constraints on the possible properties of (e) Majorana neutrinos.** |
| **(10) Double-beta decay is (=) a nuclear process in which two neutrons turn into (e&eb) two protons (or vice versa), emitting (eb) two antineutrinos.** |
| **(11) If neutrinos are Majorana fermions, (eb) the two antineutrinos could annihilate each other, and some double-beta decays would emit (eb) no antineutrinos.** |
| **(12) However, neutrinoless decay is believed to be (=) extremely rare. Searches of the process using (e) large amounts of the 35 naturally occurring (eb) isotopes that can undergo (e&eb) double-beta decay (such as xenon-136) have so far come up empty.** |

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| **(13) The KamLAND-Zen collaboration boosted their experimental sensitivity to (e) neutrinoless double-beta decays through (e&eb) a combination of factors.** |
| **(14) The researchers deployed clean detectors with (e&eb) very low background noise.** |
| **(15) And they used an unprecedented amount of xenon-136, which they purified to (e) remove any radioactive contaminants that could produce (eb) unwanted signals in the detectors.** |
| **(16) Combining data from several years, they improved (eb) the limit on the probability of neutrinoless double-beta decay by (e) sixfold compared to previous searches.** |
| **(17) This probability corresponds to (e&eb) a decay lifetime whose staggering value exceeds (e) 1026 years.** |
| **(18) This measurement allowed (eb) the researchers to set (eb) the most stringent upper limit on (e&eb) the possible mass of Majorana neutrinos (the particle must be lighter than 61–165 meV). This research is published in Physical Review Letters. –Matteo Rini Matteo Rini is the Deputy Editor of Physics.** |
| **How to Reach a Thousand-Second in-Plane Polarization Lifetime with 0.97−GeV/c Deuterons in a Storage Ring G. Guidoboni et al. (JEDI Collaboration) Phys. Rev. Lett. 117, 054801 – Published 28 July 2016 Physics logo See Synopsis: Deuterons Spin together for a Thousand Seconds** |
| **(19) Authors observe a deuteron beam polarization lifetime near 1000 s in (eb) the horizontal plane of (e) a magnetic storage ring (COSY).** |
| **(20) This long spin coherence time is maintained through (e&eb) a combination of beam bunching, electron cooling, (e&eb) sextupole field corrections, and (e&eb) the suppression of collective effects through (e&eb) beam current limits.** |
| **(21) This record lifetime is required for (e) a storage ring search for (e) an intrinsic electric dipole moment on (e&eb) the deuteron at a statistical sensitivity level approaching 10−29 e cm.** |
| **Synopsis: Deuterons Spin together for a Thousand Seconds July 28, 2016 Researchers set a new record for the in-plane spin-alignment lifetime of deuterons circulating in a magnetic storage ring. Synopsis figure E. Stephenson/Indiana University Center for Spacetime Symmetries Sime sentences and figures are deleted due to spatial restrictions.** |
| **(22) The electric dipole moment (EDM) of (e) a fundamental particle or a simple nucleus is believed to be (=) extremely small.** |
| **(23) One might therefore think that it isn't worth measuring this quantity. But it is. A nonzero EDM would signal (eb) a broken symmetry of (e) nature known as CP violation, which is required to (e) explain why the Universe has (e) more matter than antimatter.** |
| **(24) Edward Stephenson from Indiana University, Bloomington, and colleagues has now made a significant stride towards determining (eb) the EDM of the deuteron using (e) a deuteron beam circulating in (e&eb) a magnetic storage ring.** |

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| **(25) They have been able to preserve the alignment of the deuterons' spins for (e) a chunk of time that may be long enough to determine (eb) the deuteron's EDM.** |
| **(26) In experiments of this kind, the EDM can be determined by (e) measuring the rotation of (e) the initially aligned spins of (e) the beam particles as they travel around (e&eb) the ring.** |
| **(27) To obtain a measurement sufficiently accurate to explore (e&eb) CP violation predicted by (e) theories beyond the standard model, the spin alignment has to be (=) preserved for about 1000 s.** |
| **(28) But changes in particle speed—associated with (e&eb) position or direction variations among (e&eb) the particles in the beam—lead to (eb) a rapid loss of the alignment in the plane of the ring.** |
| **(29) Stephenson and co-workers have managed to retain (eb)the alignment of deuterons for the required 1000 s, which is 3 orders of magnitude longer (eb) than the previous record (obtained for an electron beam).** |
| **(30) Their success resulted from (e) capturing the beam into a tight bunch to reduce (e) position and direction spread and from (e) shaping the ring's magnetic field to(e) ensure all of the deuteron orbits around (e&eb) the ring have the same length, despite (e) direction variations. This research is published in Physical Review Letters. –Ana Lopes Ana Lopes is a Senior Editor of Physics.** |
| **Cotunneling Drag Effect in Coulomb-Coupled Quantum Dots A. J. Keller, J. S. Lim, David Sánchez, Rosa López, S. Amasha, J. A. Katine, Hadas Shtrikman, and D. Goldhaber-Gordon Phys. Rev. Lett. 117, 066602 – Published 4 August 2016 Physics logo See Synopsis: Coulomb Drag in a Double Dot** |
| **(31) In Coulomb drag, a current flowing in one conductor can induce (eb) a voltage across an adjacent conductor via (e&eb) the Coulomb interaction.** |
| **(32) The mechanisms yielding drag effects are not always understood, even though drag effects are (=) sufficiently general to be (=) seen in many low-dimensional systems.** |
| **(33) In this Letter, authors observe Coulomb drag in (eb) a Coulomb-coupled double quantum dot and, through (e&eb) both experimental and theoretical arguments, identify cotunneling as (=) essential to obtaining a correct qualitative understanding of the drag behavior.** |
| **Synopsis: Coulomb Drag in a Double Dot August 4, 2016 Electric current passing through a quantum dot can generate current in a nearby dot through a coordinated tunneling mechanism. Synopsis figureA. J. Keller et al., Phys. Rev. Lett. (2016)** |
| **(34) Electrons moving through (e&eb) a conductor can induce (eb) the movement of electrons of a nearby, but isolated, conductor.** |
| **(35) This so-called Coulomb drag is well understood in 2D layered materials, but the theory for nanowires and quantum dots does not (e) fully explain experiments.** |
| **(36) A new study of Coulomb drag in two nearby quantum dots shows (eb) that simultaneous tunneling of electrons—called cotunneling—can help explain the low-temperature behavior of (e) the system.** |

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| **(37) Electric current flows through (e&eb) a quantum dot by the tunneling of (e) electrons.** |
| **(38) When two dots are placed next to each other, the active, or lead, dot can drag current through (e&eb) the passive dot by (e) a process called sequential tunneling.** |
| **(39) In this case, an electron first tunnels onto (e&eb) the active dot, increasing (eb) the probability of tunneling onto (e&eb) the passive dot.** |
| **(40) Thus the passive dot follows (e&eb) the active dot's lead, but always one step behind.** |
| **(41) Sequential tunneling explains (eb) Coulomb drag observations at (eb) high temperature, but no studies have investigated cold temperatures where (e) quantum coherence effects may play a role.** |
| **(42) Andrew Keller from the California Institute of Technology, Pasadena, and his colleagues explored Coulomb drag at (eb) sub-Kelvin temperatures in a pair of quantum dots coupled by (e) capacitance.** |
| **(43) The team applied a current-driving voltage across (e&eb) one dot and detected a drag current in (eb) the other dot, even in regimes where (e) the driving voltage was very small.** |
| **(44) This conflicts with (e&eb) the sequential tunneling model, which predicts (eb) a threshold driving voltage below which (e) the drag current should be zero.** |
| **(45) The team showed that the inclusion of cotunneling into (eb) the drag theory explains the observations.** |
| **(46) Cotunneling could be important to efforts aimed at harvesting a unidirectional current from (e) a fluctuating voltage supply. This research is published in Physical Review Letters. –Michael Schirber Michael Schirber is a Corresponding Editor for Physics based in Lyon, France.** |
| **Maybe it wasn't the Higgs particle after all November 7, 2014 Maybe it wasn't the Higgs particle after all Associate Professor Mads Toudal Frandsen, University of Southern Denmark Credit: University of Southern Denmark** |
| **(47) Last year CERN announced the finding of a new elementary particle, the Higgs particle. But maybe it wasn't the Higgs particle, maybe it just looks like it. And maybe it is not alone. Many calculations indicate that the particle discovered last year in the CERN particle accelerator was (=) indeed the famous Higgs particle.** |
| **(48) Physicists agree that the CERN experiments did find a new particle that had never been seen before, but according to (e&eb) an international research team, there is (=) no conclusive evidence that the particle was indeed the Higgs particle.The research team has scrutinized the existing scientific data from CERN about the newfound particle and published their analysis in the journal Physical Review D. A member of this team is Mads Toudal Frandsen, associate professor at the Center for Cosmology and Particle Physics Phenomenology, Department of Physics, Chemistry and Pharmacy at the University of Southern Denmark. "The CERN data is generally taken as (=) evidence that the particle is the Higgs particle.** |

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| **(49) It is true that the Higgs particle can explain (eb) the data but there can be other explanations, we would also get this data from other particles", Mads Toudal Frandsen explains.** |
| **(50) The researchers' analysis does not (e) debunk the possibility that CERN has discovered the Higgs particle. That is still possible - but it is equally possible that it is a different kind of particle. "The current data is not precise enough to determine exactly what the particle is. It could be a number of other known particles", says Mads Toudal Frandsen. But if it wasn't the Higgs particle, which was found in CERN's particle accelerator, then what was it? "We believe that it may be a so-called techni-higgs particle. This particle is in some ways similar to the Higgs particle - hence half of the name", says Mads Toudal Frandsen.** |
| **(51) Although the techni-higgs particle and Higgs particle can easily be confused in experiments, they are two very different particles belonging to (e) two very different theories of how the universe was created.** |
| **(52) The Higgs particle is the missing piece in (eb) the theory called the Standard Model.** |
| **(53) This theory describes (eb) three of the four forces of nature. But it does not (e) explain what dark matter is - the substance that makes up (eb) most of the universe. A techni-higgs particle, if it exists, is (=) a completely different thing:** |
| **(54) "A techni-higgs particle is not (e) an elementary particle.** |
| **(55) Instead, it consists of so-called techni-quarks, which we believe are elementary. Techni-quarks may bind together in various ways to form (eb) for instance techni-higgs particles, while other combinations may form (eb) dark matter. Authors therefore expect to find several different particles at the LHC, all built by (e) techni-quarks", says Mads Toudal Frandsen.** |
| **(56) If techni-quarks exist, (eb) there must be a force to bind them together so that (e) they can form particles.** |
| **(57) None of the four known forces of nature (gravity, the electromagnetic force, the weak nuclear force and the strong nuclear force) are (=) any good at binding techni-quarks together.** |
| **(58) There must therefore be (=) a yet undiscovered force of nature. This force is called the the Technicolor force.** |
| **(59) What was found last year in CERN's accelerator could thus be (=) either the Higgs particle of the Standard Model or a light techni-higgs particle, composed of (e) two techni-quarks.** |
| **(60) Mads Toudal Frandsen believes that more data from CERN will probably be able to determine (eb) if it was a Higgs or a techni-higgs particle. If CERN gets an even more powerful accelerator, it will in principle be able to observe techni-quarks directly. More information: Technicolor Higgs boson in the light of LHC data. Phys. Rev. D 90, 035012th Alexander Belyaev, Matthew S. Brown, Roshan Foadi, and Mads T. Frandsen. journals.aps.org/prd/abstract/10.1103/PhysRevD.90.035012. On Arxiv: arxiv.org/abs/1309.2097 Journal reference: Physical Review D search and more info website: Provided by: University of Southern Denmark search and more info website** |

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| **The inflationary brane-antibrane universeClifford P. Burgess1,2, Mahbub Majumdar3, Detlef Nolte1, Fernando Quevedo3, Govindan Rajesh1 and Ren-Jie Zhang1 Published 13 August 2001 • Journal of High Energy Physics, Volume 2001, JHEP07(2001)** |
| **(61) Authors show how the motion through (e&eb) the extra dimensions of a gas of branes and antibranes can, under certain circumstances, produce (eb) an era of inflation as seen by (e) observers trapped on (e&eb) a 3-brane, with (e&eb) the inflaton being the inter-brane separation.** |
| **(62) Although most of discussion refers to (e) arbitrary ps, when we need to be specific they assume that they are (=) D-branes of type-II or type-I string theory.** |
| **(63) For realistic brane couplings, such as those arising in (eb) string theory, the inter-brane potentials are (=) too steep to inflate (eb) the universe for acceptably long times.** |
| **(64) However, for special regions of the parameter space of (e) brane-antibrane positions the brane motion is (=) slow enough for there to be (=) sufficient inflation.** |
| **(65) Inflation would be more generic in models where (e) the inter-brane interactions are much weaker.** |
| **(66) The spectrum of primordial density fluctuations predicted has (e) index n slightly less than 1, and acceptable amplitude, provided (e) that the extra dimensions have (e) linear size 1/r ~ 1012 GeV.** |
| **(67) Reheating occurs as in hybrid inflation, with (e&eb) the tachyonic instability of (e) the brane-antibrane system taking over for (e) small separations.** |
| **(68) The tachyon field can induce (eb) a cascade mechanism within which higher-dimension branes annihilate (e) into lower-dimension ones.** |
| **(69) They argue that such a cascade naturally stops with (e&eb) the production of 3-branes in (eb) 10-dimensional string theory.** |
| **Effect of the neutron lifetime on processes in the early universe A. V. Chechkin, A. V. Ivanchik, A. P. Serebrov, S. V. Bobashev First Online: 27 July 2016 DOI: 10.1134/S1063784216070069 Chechkin, A.V., Ivanchik, A.V., Serebrov, A.P. et al. Tech. Phys. (2016) 61: 1101. Doi: 10.1134/S1063784216070069** |
| **(70) The influence of the neutron lifetime on (e&eb) the abundance of light elements produced during (e&eb) the primordial nucleosynthesis early in (eb) the birth of the Universe is considered.** |
| **(71) Among light elements, namely, D, 3He, 4He, and 7Li, 4He proves to be (=) most sensitive to neutron lifetime τn.** |
| **(72) Astronomic data on the light element abundance also provide (eb) the best accuracy for 4He.** |
| **(73) The solution of a number of problems discussed in this paper requires (e) improving the accuracy of observations for the 4He abundance and refining the value of τn.** |

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| **Application of a Theory and Simulation based Convective Boundary Mixing model for AGB Star Evolution and Nucleosynthesis U. Battino, M. Pignatari, C. Ritter, F. Herwig, P. Denisenkov, J.W. Den Hartogh, R. Trappitsch, R. Hirschi, B. Freytag, F. Thielemann, B. Paxton (Submitted on 19 May 2016)** |
| **(74) The s-process nucleosynthesis in Asymptotic Giant Branch (AGB) stars depends on (e&eb) the modeling of (e) convective boundaries.** |
| **(75) Authors' present models and s-process simulations that adopt (e&eb) a treatment of convective boundaries based on (e) the results of hydrodynamic simulations and on (e&eb) the theory of mixing due to (e) gravity waves in (eb) the vicinity of convective boundaries.** |
| **(76) Hydrodynamics simulations suggest (eb) the presence of convective boundary mixing (CBM) at (eb) the bottom of the thermal pulse-driven convective zone.** |
| **(77) Similarly, convection-induced mixing processes are proposed for (e) the mixing below the convective envelope during (e&eb) third dredge-up where (e) the 13C pocket for (e) the s process in AGB stars forms.** |
| **(78) In this work apply apply a CBM model motivated by (e) simulations and theory to (e&eb) models with initial mass M = 2 and M = 3M?, and with (e&eb) initial metal content Z = 0:01 and Z = 0:02.** |
| **(79) As reported previously, the He-intershell abundance of 12C and 16O are increased by (e) CBM at the bottom of (e) pulse-driven convection zone.** |
| **(80) This mixing is affecting (e&eb) the 22Ne (alpha, n) 25Mg activation and (e&eb) the s-process effciency in the 13C-pocket.** |
| **(81) In our model CBM at the bottom of the convective envelope during (e&eb) the third dredgeup represents (eb) gravity wave mixing.** |
| **(82) Authors take further into account that hydrodynamic simulations indicate (eb) a declining mixing efficiency already about (e&eb) a pressure scale height from (e) the convective boundaries, compared to (e&eb) mixing-length theory.** |
| **(83) The final s-process abundances are characterized by (e) 0:36 < [s=Fe] < 0:78 and the heavy-to-light s-process ratio is (=) 0:23 < [hs=ls] < 0:45.** |
| **(84) Finally, authors compare results with (e&eb) stellar observations, pre-solar grain measurements and previous work. Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI:10.3847/0004-637X/827/1/30 Cite as: arXiv:1605.06159 [astro-ph.SR] (or arXiv:1605.06159v1 [astro-ph.SR] for this version)** |
| **MESA and NuGrid simulations of classical novae: CO and ONe nova nucleosynthesis Pavel A. Denissenkov (UVic), James W. Truran (Chicago), Marco Pignatari (Basel), Reto Trappitsch (Chicago), Christian Ritter (Frankfurt), Falk Herwig (UVic), Umberto Battino (Basel), Kiana Setoodehnia (McMaster), B. Paxton (UC Santa Barbara)** |

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| **(85) Classical novae are the result of (e) thermonuclear flashes of hydrogen accreted by (e) CO or ONe white dwarfs, leading (eb) eventually to the dynamic ejection of (e) the surface layers.** |
| **(86) These are observationally known to be (=) enriched in heavy elements, such as (=) C, O and Ne that must originate in (eb) layers below the H-flash convection zone.** |
| **(87) Building on their previous work, they now present stellar evolution simulations of (e) ONe novae and provide (eb) a comprehensive comparison of models with (e&eb) published ones.** |
| **(88) Some of authors models include exponential convective boundary mixing to account for (e) the observed enrichment of (e) the nova ejecta even when accreted material has (e) a solar abundance distribution.** |
| **(89) Models produce (eb) maximum temperature evolution profiles and nucleosynthesis yields (eb) in good agreement with (e&eb) models that generate (eb) enriched ejecta by assuming (eb) that the accreted material was pre-mixed.** |
| **(90) Authors confirm for (e) ONe novae the result we reported previously, i.e.\ we found that 3He could be produced {\it in situ} in (eb) solar-composition envelopes accreted with (e&eb) slow rates ($M˙$<10−10M⊙/yr) by (e) cold (TWD<107 K) CO WDs, and that convection was triggered by (eb) 3He burning before (e) the nova outburst in that case.** |
| **(91) In addition, authors now find that the interplay between the 3He production and (e&eb) destruction in the solar-composition envelope accreted with (e&eb) an intermediate rate, e.g.\ $M˙$=10−10M⊙/yr, by the 1.15M⊙ ONe WD with (e&eb) a relatively high initial central temperature, e.g.\ TWD=15×106 K, leads to (eb) the formation of a thick radiative buffer zone that separates (e&eb) the bottom of the convective envelope from (e) the WD surface. (Abridged) Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI: 10.1093/mnras/stu1000 Cite as: arXiv:1303.6265 [astro-ph.SR] (or arXiv:1303.6265v2 [astro-ph.SR] for this version)** |
| **Hybrid C-O-Ne White Dwarfs as Progenitors of Diverse SNe Ia Pavel Denissenkov (UVic), James Truran (U. Chicago), Falk Herwig (UVic), Sam Jones (UVic), Bill Paxton (UC Santa Barbara), Ken'ichi Nomoto (U. Tokyo), Toshio Suzuki (Nihon U.), Hiroshi Toki (Osaka U.)(Submitted on 6 Nov 2014)** |
| **(92) When carbon is ignited off-center in a CO core of a super-AGB star (eb) its burning in a convective shell tends to propagate to (e&eb) the center.** |
| **(93) Whether the C flame will actually be able to reach (eb) the center depends on (e&eb) the efficiency of (e) extra mixing beneath (e&eb) the C convective shell.** |
| **(94) Whereas thermohaline mixing is too inefficient to interfere with the C-flame propagation, (eb) convective boundary mixing can prevent (e) the C burning from reaching (eb) the center.** |
| **(95) As a result, a C-O-Ne white dwarf (WD) is (=) formed, after the star has lost (e) its envelope.** |
| **(96) Such a "hybrid" WD has (e) a small CO core surrounded by (e) a thick ONe zone.** |

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| **(97) In the proposed 1D stellar evolution computations the hybrid WD is allowed to (eb) accrete C-rich material, as if it were in a close binary system and (e&eb) accreted H-rich material from (e) its companion with (e&eb) a sufficiently high rate at which the accreted H would be processed into (e&eb) He under (e&eb) stationary conditions, assuming (eb) that He could then be transformed into (e&eb) C.** |
| **(98) When the mass of the accreting WD approaches (eb) the Chandrasekhar limit, they find a series of convective Urca shell flashes associated with (e&eb) high abundances of 23Na and 25Mg.** |
| **(99) They are followed by off-center C ignition leading to (eb) convection that occupies (e) almost the entire star.** |
| **(100) To model the Urca processes, authors use the most recent well-resolved data for (e) their reaction and (e&eb) neutrino-energy loss rates.** |
| **(101) Because of the emphasized uncertainty of the convective Urca process in hybrid WD models of (e) SN Ia progenitors, authors consider a number of their potentially possible alternative instances for (e) different mixing assumptions, all of which reach (eb) a phase of explosive C ignition, either off or in (eb) the center.** |
| **(102) Our hybrid SN Ia progenitor models have (e) much lower C to O abundance ratios at (eb) the moment of the explosive C ignition than (e) their pure CO counterparts, which may explain (eb) the observed diversity of the SNe Ia. subjects: Solar and Stellar Astrophysics (astro-ph.SR)cite as: arXiv:1411.1471 [astro-ph.SR](or arXiv:1411.1471v1 [astro-ph.SR] for this version)** |
| **Rejecting proposed dense-matter equations of state with quiescent low-mass X-ray binaries Sebastien Guillot, Robert E. Rutledge (Submitted on 15 Sep 2014 (v1), last revised 20 Oct 2014 (this version, v2))** |
| **(103) Neutrons stars are unique laboratories to (e) discriminate between the various proposed equations of (e) state of matter at and above nuclear density.** |
| **(104) One sub-class of neutron stars - those inside (eb) quiescent low-mass X-ray binaries (qLMXBs) - produce (eb) a thermal surface emission from (e) which the neutron star radius (R_NS) can be measured, using (e) the widely accepted observational scenario for qLMXBs, assuming (eb) unmagnetized H atmospheres.** |
| **(105) In a combined spectral analysis, this work first reproduces (eb) a previously published measurement of the \rns, assumed to be (=) the same for all neutron stars, using (e) a slightly expanded data set.** |
| **(106) The radius measured is R_NS = 9.4 +/-1.2 km. On the basis of spectral analysis alone, this measured value is not affected by imposing (e&eb) an assumption of causality in the core.** |
| **(107) However, the assumptions underlying this R_NS measurement would be falsified by (e) the observation of any neutron star with (e&eb) a mass >2.6 Msun, since radii <11 km would be rejected if (e) causality is assumed, which would exclude (eb) most of the R_NS parameter space obtained in this analysis.** |

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| **(108) Finally, this work directly tests (e&eb) a selection of dense matter equations of states: WFF1, AP4, MPA1, PAL1, MS0, and three versions of equations of state produced through (e&eb) chiral effective theory.** |
| **(109) Two of those, MSO and PAL1, are rejected at (eb) the 99% confidence level, accounting for (e) all quantifiable uncertainties, while (e) the other cannot be excluded at >99% certainty. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE) DOI: 10.1088/2041-8205/796/1/L3 Cite as: arXiv:1409.4306 [astro-ph.HE] (or arXiv:1409.4306v2 [astro-ph.HE] for this version** |

## SECTION THREE

## Recoil-Heating Rate, Which Depends On The Photon-Scattering Rate And The Effect Of Each Scattering Event

**INTRODUCTION—VARIABLES USED**

**Viewpoint: Measuring Quantum Kicks from a Beam of Light Amir Safavi-Naeini, 348 Via Pueblo Mall, Stanford University, Stanford, CA 94305-4088, USA June 13, 2016• Physics 9, 64 Force sensors levitated by light have reached the quantum regime, in which their sensitivity is limited by the momentum kicks of individual photons. Figure caption expand figure APS/Alan Stonebraker Figures and some sentences are deleted due to spatial restrictions Please excuse me**

(1) The time scale needed for (e) such a change is determined by (e) the recoil-heating rate, which depends on (e&eb) the photon-scattering rate and the effect of (e) each scattering event.
(2) The achievement of Novotny's group has been to measure (eb) the recoil-heating rate for (e) an object as (=) large as a 50-nm-diameter nanosphere.
(3) For this measurement to succeed, the team had to reduce (e) the motional energy of the sphere to (e) an extremely low level, corresponding to (e&eb) occupancy of 60 phonons in (eb) the harmonic trap, or a temperature of (e) a few hundred microkelvins.
(4) Such cooling is typically accomplished by applying radiation pressure from (e) a laser field along (e&eb) three orthogonal positions in the trap independently.
(5) In contrast, the ETH team utilized (e) a parametric cooling process.
(6) This involves reading out the particle's motion in (eb) each direction and then modulating (e&eb) the intensity of the trapping laser so as to adjust (e&eb) a parameter (in this case, the natural frequency of (e) the harmonic trap) that deamplifies (e) the particle's motion [8].
(7) The process is a bit like standing on (e&eb) a swing to decrease its amplitude.
(8) Under optimal conditions, the researchers obtained (eb) a balance between cooling from the lasers and heating from noise that leads to (eb) an equilibrium temperature of (e) about 500 $\mu\mu$K.

**NOTATION**

**Module One**

However, the cumulative momentum kicks from many scattering events can change (e&eb) the state of the nanosphere

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| $G_{13}$ : Category one of **cumulative momentum kicks from many scattering events;** the state of the nanosphere<br><br>$G_{14}$ : Category two of SAS<br><br>$G_{15}$ : Category three of SAS | |
| $T_{13}$ : Category one of the state of the nanosphere; **cumulative momentum kicks from many scattering events**<br><br>$T_{14}$ : Category two of SAS<br><br>$T_{15}$ : Category three of SAS | |
| **Module Two**<br><br>The time scale needed for (e) such a change is determined by (e) the recoil-heating rate, which depends on (e&eb) the photon-scattering rate and the effect of (e) each scattering event | |
| $G_{16}$ : Category one of such a change is determined by (e) the recoil-heating rate, which depends on (e&eb) the photon-scattering rate and the effect of (e) each scattering event<br><br>$G_{17}$ : Category two of SAS<br><br>$G_{18}$ : Category three of SAS | |
| $T_{16}$ : Category one of time scale needed<br><br>$T_{17}$ : Category two of SAS<br><br>$T_{18}$ : Category three of SAS | |
| **Module three**<br><br>The time scale needed for such a change is determined by (e) the recoil-heating rate, which depends on (e&eb) the photon-scattering rate and the effect of (e) each scattering event | |
| $G_{20}$ : Category one of<br><br>$G_{21}$ : Category two of SAS<br><br>$G_{22}$ : Category three of SAS | |
| $T_{20}$ : Category one of time scale needed for such a change is determined<br><br>$T_{21}$ : Category two of SAS<br><br>$T_{22}$ : Category three of SAS | |
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| recoil-heating rate, which depends on the photon-scattering rate and the effect of (e) each scattering event | |
| $G_{24}$ : Category one of each scattering<br>$G_{25}$ : Category two of SAS<br>$G_{26}$ : Category three of SAS | |
| $T_{24}$ : Category one of recoil-heating rate, which depends on the photon-scattering rate and the effect<br>$T_{25}$ : Category two of SAS<br>$T_{26}$ : Category three of SAS | |
| **Module five**<br>For this measurement to succeed, the team had to reduce (e) the motional energy of the sphere to (e) an extremely low level, corresponding to (e&eb) occupancy of 60 phonons in (eb) the harmonic trap, or a temperature of (e) a few hundred microkelvins | |
| $G_{28}$ : Category one of measurement to succeed<br>$G_{29}$ : Category two of SAS<br>$G_{30}$ : Category three of SAS | |
| $T_{28}$ : Category one ofteam had to reduce (e) the motional energy of the sphere to (e) an extremely low level, corresponding to (e&eb) occupancy of 60 phonons in (eb) the harmonic trap, or a temperature of (e) a few hundred microkelvins<br>$T_{29}$ : Category two of SAS<br>$T_{30}$ : Category three of SAS | |
| **Module six**<br>For this measurement to succeed, the team had to reduce (e) the motional energy of the sphere to (e) an extremely low level, corresponding to (e&eb) occupancy of 60 phonons in (eb) the harmonic trap, or a temperature of (e) a few hundred microkelvins | |
| $G_{32}$ : Category one of **measurement to succeed, the team had to reduce (e) the motional energy of the sphere to (e) an extremely low level, corresponding;** occupancy of 60 phonons in (eb) the harmonic trap, or a temperature of (e) a few hundred microkelvins<br>$G_{33}$ : Category two of SAS<br>$G_{34}$ : Category three of SAS | |
| $T_{32}$ : Category one of occupancy of 60 phonons in (eb) the harmonic trap, or a temperature of (e) a few hundred microkelvins **;measurement to succeed, the team had to reduce (e) the motional energy of the sphere to (e) an extremely low level, corresponding** | |

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| $T_{33}$ : Category two of SAS<br><br>$T_{34}$ : Category three of SAS | |
| **Module seven**<br><br>For this measurement to succeed, the team had to reduce (e) the motional energy of the sphere to (e) an extremely low level, corresponding to (e&eb) occupancy of 60 phonons in (eb) the harmonic trap, or a temperature of (e) a few hundred microkelvins | |
| $G_{36}$ : Category one of motional energy of the sphere to (e) an extremely low level, corresponding to (e&eb) occupancy of 60 phonons in (eb) the harmonic trap, or a temperature of (e) a few hundred microkelvins<br><br>$G_{37}$ : Category two of SAS<br><br>$G_{38}$ : Category three of SAS | |
| $T_{36}$ : Category one of measurement to succeed, the team had to<br><br>$T_{37}$ : Category two of SAS<br><br>$T_{38}$ : Category three of SAS | |
| **Module eight**<br><br>For this measurement to succeed, the team had to reduce (e) the motional energy of the sphere to an extremely low level, corresponding to (e&eb) occupancy of 60 phonons in (eb) the harmonic trap, or a temperature of (e) a few hundred microkelvins | |
| $G_{40}$ : Category one of **measurement to succeed, the team had to reduce (e) the motional energy of the sphere to an extremely low level;** occupancy of 60 phonons in (eb) the harmonic trap, or a temperature of (e) a few hundred microkelvins<br><br>$G_{41}$ : Category two of SAS<br><br>$G_{42}$ : Category three of SAS | |
| $T_{40}$ : Category one of occupancy of 60 phonons in (eb) the harmonic trap, or a temperature of (e) a few hundred microkelvins **measurement to succeed, the team had to reduce (e) the motional energy of the sphere to an extremely low level**<br><br>$T_{41}$ : Category two of SAS<br><br>$T_{42}$ : Category three of SAS | |
| **Module Nine**<br><br>For this measurement to succeed, the team had to reduce the motional energy of the sphere to an extremely low level, corresponding to (e&eb) occupancy of 60 phonons in (eb) the harmonic trap, or a temperature of (e) a few hundred microkelvins | |

$G_{44}$ : Category one of **measurement to succeed, the team had to reduce the motional energy of the sphere to an extremely low level;** occupancy of 60 phonons in (eb) the harmonic trap, or a temperature of (e) a few hundred microkelvins

$G_{45}$ : Category two of SAS

$G_{46}$ : Category three of SAS

$T_{44}$ : Category one of occupancy of 60 phonons in (eb) the harmonic trap, or a temperature of (e) a few hundred microkelvins; **measurement to succeed, the team had to reduce the motional energy of the sphere to an extremely low level,**

$T_{45}$ : Category two of SAS

$T_{46}$ : Category three of SAS

**The Coefficients:**

$(a_{13})^{(1)}, (a_{14})^{(1)}, (a_{15})^{(1)}, (b_{13})^{(1)}, (b_{14})^{(1)}, (b_{15})^{(1)}\ (a_{16})^{(2)}, (a_{17})^{(2)}, (a_{18})^{(2)}\ (b_{16})^{(2)}, (b_{17})^{(2)}, (b_{18})^{(2)}$:
$(a_{20})^{(3)}, (a_{21})^{(3)}, (a_{22})^{(3)}$ , $(b_{20})^{(3)}, (b_{21})^{(3)}, (b_{22})^{(3)}$
$(a_{24})^{(4)}, (a_{25})^{(4)}, (a_{26})^{(4)}, (b_{24})^{(4)}, (b_{25})^{(4)}, (b_{26})^{(4)}, (b_{28})^{(5)}, (b_{29})^{(5)}, (b_{30})^{(5)}, (a_{28})^{(5)}, (a_{29})^{(5)}, (a_{30})^{(5)}$,
$(a_{32})^{(6)}, (a_{33})^{(6)}, (a_{34})^{(6)}, (b_{32})^{(6)}, (b_{33})^{(6)}, (b_{34})^{(6)}$
$(a_{36})^{(7)}, (a_{37})^{(7)}, (a_{38})^{(7)}, (b_{36})^{(7)}, (b_{37})^{(7)}, (b_{38})^{(7)}$
$(a_{40})^{(8)}, (a_{41})^{(8)}, (a_{42})^{(8)}, (b_{40})^{(8)}, (b_{41})^{(8)}, (b_{42})^{(8)}$
$(a_{44})^{(9)}, (a_{45})^{(9)}, (a_{46})^{(9)}, (b_{44})^{(9)}, (b_{45})^{(9)}, (b_{46})^{(9)}$

are Accentuation coefficients
$(a'_{13})^{(1)}, (a'_{14})^{(1)}, (a'_{15})^{(1)}, (b'_{13})^{(1)}, (b'_{14})^{(1)}, (b'_{15})^{(1)}, (a'_{16})^{(2)}, (a'_{17})^{(2)}, (a'_{18})^{(2)}$, $(b'_{16})^{(2)}, (b'_{17})^{(2)}, (b'_{18})^{(2)}$
$, (a'_{20})^{(3)}, (a'_{21})^{(3)}, (a'_{22})^{(3)}, (b'_{20})^{(3)}, (b'_{21})^{(3)}, (b'_{22})^{(3)}$
$(a'_{24})^{(4)}, (a'_{25})^{(4)}, (a'_{26})^{(4)}, (b'_{24})^{(4)}, (b'_{25})^{(4)}, (b'_{26})^{(4)}, (b'_{28})^{(5)}, (b'_{29})^{(5)}, (b'_{30})^{(5)}\ (a'_{28})^{(5)}, (a'_{29})^{(5)}, (a'_{30})^{(5)}$ ,
$(a'_{32})^{(6)}, (a'_{33})^{(6)}, (a'_{34})^{(6)}, (b'_{32})^{(6)}, (b'_{33})^{(6)}, (b'_{34})^{(6)}$
$(a'_{36})^{(7)}, (a'_{37})^{(7)}, (a'_{38})^{(7)}, (b'_{36})^{(7)}, (b'_{37})^{(7)}, (b'_{38})^{(7)}$,
$(a'_{40})^{(8)}, (a'_{41})^{(8)}, (a'_{42})^{(8)}, (b'_{40})^{(8)}, (b'_{41})^{(8)}, (b'_{42})^{(8)}$,
$(a'_{44})^{(9)}, (a'_{45})^{(9)}, (a'_{46})^{(9)}, (b'_{44})^{(9)}, (b'_{45})^{(9)}, (b'_{46})^{(9)}$,
are Dissipation coefficients

**Module Numbered One**

**The differential system of this model is now (Module Numbered one)**

$$\frac{dG_{13}}{dt} = (a_{13})^{(1)}G_{14} - \left[(a'_{13})^{(1)} + (a''_{13})^{(1)}(T_{14}, t)\right]G_{13} \quad (1)$$

$$\frac{dG_{14}}{dt} = (a_{14})^{(1)}G_{13} - \left[(a'_{14})^{(1)} + (a''_{14})^{(1)}(T_{14}, t)\right]G_{14} \quad (2)$$

$$\frac{dG_{15}}{dt} = (a_{15})^{(1)}G_{14} - \left[(a'_{15})^{(1)} + (a''_{15})^{(1)}(T_{14}, t)\right]G_{15} \quad (3)$$

$$\frac{dT_{13}}{dt} = (b_{13})^{(1)}T_{14} - \left[(b'_{13})^{(1)} - (b''_{13})^{(1)}(G, t)\right]T_{13} \quad (4)$$

$$\frac{dT_{14}}{dt} = (b_{14})^{(1)}T_{13} - \left[(b'_{14})^{(1)} - (b''_{14})^{(1)}(G, t)\right]T_{14} \quad (5)$$

$$\frac{dT_{15}}{dt} = (b_{15})^{(1)}T_{14} - \left[(b'_{15})^{(1)} - (b''_{15})^{(1)}(G, t)\right]T_{15} \quad (6)$$

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| $+(a_{13}^{''})^{(1)}(T_{14},t) =$ First augmentation factor | |
| $-(b_{13}^{''})^{(1)}(G,t) =$ First detritions factor | |
| **Module Numbered Two**<br><br>**The differential system of this model is now ( Module numbered two)** | |
| $\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \left[(a_{16}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17},t)\right]G_{16}$ | 7 |
| $\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \left[(a_{17}^{'})^{(2)} + (a_{17}^{''})^{(2)}(T_{17},t)\right]G_{17}$ | 8 |
| $\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \left[(a_{18}^{'})^{(2)} + (a_{18}^{''})^{(2)}(T_{17},t)\right]G_{18}$ | 9 |
| $\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \left[(b_{16}^{'})^{(2)} - (b_{16}^{''})^{(2)}\left((G_{19}),t\right)\right]T_{16}$ | 10 |
| $\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \left[(b_{17}^{'})^{(2)} - (b_{17}^{''})^{(2)}\left((G_{19}),t\right)\right]T_{17}$ | 11 |
| $\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \left[(b_{18}^{'})^{(2)} - (b_{18}^{''})^{(2)}\left((G_{19}),t\right)\right]T_{18}$ | 12 |
| $+(a_{16}^{''})^{(2)}(T_{17},t) =$ First augmentation factor | |
| $-(b_{16}^{''})^{(2)}\left((G_{19}),t\right) =$ First detritions factor | |
| **Module Numbered Three**<br><br>**The differential system of this model is now (Module numbered three)** | |
| $\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \left[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21},t)\right]G_{20}$ | 13 |
| $\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \left[(a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)}(T_{21},t)\right]G_{21}$ | 14 |
| $\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \left[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21},t)\right]G_{22}$ | 15 |
| $\frac{dT_{20}}{dt} = (b_{20})^{(3)}T_{21} - \left[(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23},t)\right]T_{20}$ | 16 |
| $\frac{dT_{21}}{dt} = (b_{21})^{(3)}T_{20} - \left[(b_{21}^{'})^{(3)} - (b_{21}^{''})^{(3)}(G_{23},t)\right]T_{21}$ | 17 |
| $\frac{dT_{22}}{dt} = (b_{22})^{(3)}T_{21} - \left[(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23},t)\right]T_{22}$ | 18 |
| $+(a_{20}^{''})^{(3)}(T_{21},t) =$ First augmentation factor | |
| $-(b_{20}^{''})^{(3)}(G_{23},t) =$ First detritions factor | |
| **Module Numbered Four**<br><br>**The differential system of this model is now (Module numbered Four)** | |
| $\frac{dG_{24}}{dt} = (a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25},t)\right]G_{24}$ | 19 |
| $\frac{dG_{25}}{dt} = (a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25},t)\right]G_{25}$ | 20 |
| $\frac{dG_{26}}{dt} = (a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25},t)\right]G_{26}$ | 21 |
| $\frac{dT_{24}}{dt} = (b_{24})^{(4)}T_{25} - \left[(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}\left((G_{27}),t\right)\right]T_{24}$ | 22 |
| $\frac{dT_{25}}{dt} = (b_{25})^{(4)}T_{24} - \left[(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}\left((G_{27}),t\right)\right]T_{25}$ | 23 |
| $\frac{dT_{26}}{dt} = (b_{26})^{(4)}T_{25} - \left[(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}\left((G_{27}),t\right)\right]T_{26}$ | 24 |
| $+(a_{24}^{''})^{(4)}(T_{25},t) =$ **First augmentation factor** | |
| $-(b_{24}^{''})^{(4)}\left((G_{27}),t\right) =$ **First detritions factor** | |
| **Module Numbered Five**: | |

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| **The differential system of this model is now (Module number five)** | |
| $\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29},t)\right]G_{28}$ | 25 |
| $\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29},t)\right]G_{29}$ | 26 |
| $\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \left[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29},t)\right]G_{30}$ | 27 |
| $\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \left[(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}((G_{31}),t)\right]T_{28}$ | 28 |
| $\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \left[(b_{29}^{'})^{(5)} - (b_{29}^{''})^{(5)}((G_{31}),t)\right]T_{29}$ | 29 |
| $\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \left[(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}((G_{31}),t)\right]T_{30}$ | 30 |
| $+(a_{28}^{''})^{(5)}(T_{29},t) =$ **First augmentation factor** | |
| $-(b_{28}^{''})^{(5)}((G_{31}),t) =$ **First detritions factor** | |
| **Module Numbered Six**<br><br>**The differential system of this model is now (Module numbered Six)** | |
| $\frac{dG_{32}}{dt} = (a_{32})^{(6)}G_{33} - \left[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33},t)\right]G_{32}$ | 31 |
| $\frac{dG_{33}}{dt} = (a_{33})^{(6)}G_{32} - \left[(a_{33}^{'})^{(6)} + (a_{33}^{''})^{(6)}(T_{33},t)\right]G_{33}$ | 32 |
| $\frac{dG_{34}}{dt} = (a_{34})^{(6)}G_{33} - \left[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33},t)\right]G_{34}$ | 33 |
| $\frac{dT_{32}}{dt} = (b_{32})^{(6)}T_{33} - \left[(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}((G_{35}),t)\right]T_{32}$ | 34 |
| $\frac{dT_{33}}{dt} = (b_{33})^{(6)}T_{32} - \left[(b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}((G_{35}),t)\right]T_{33}$ | 35 |
| $\frac{dT_{34}}{dt} = (b_{34})^{(6)}T_{33} - \left[(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}((G_{35}),t)\right]T_{34}$ | 36 |
| $+(a_{32}^{''})^{(6)}(T_{33},t) =$ **First augmentation factor** | |
| **Module Numbered Seven:**<br><br>**The differential system of this model is now (Seventh Module)** | |
| $\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \left[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37},t)\right]G_{36}$ | 37 |
| $\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \left[(a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37},t)\right]G_{37}$ | *38* |
| $\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \left[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37},t)\right]G_{38}$ | *39* |
| $\frac{dT_{36}}{dt} = (b_{36})^{(7)}T_{37} - \left[(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}((G_{39}),t)\right]T_{36}$ | *40* |
| $\frac{dT_{37}}{dt} = (b_{37})^{(7)}T_{36} - \left[(b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}((G_{39}),t)\right]T_{37}$ | *41* |
| $\frac{dT_{38}}{dt} = (b_{38})^{(7)}T_{37} - \left[(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}((G_{39}),t)\right]T_{38}$ | *42* |
| $+(a_{36}^{''})^{(7)}(T_{37},t) =$ **First augmentation factor** | |
| **Module Numbered Eight**<br><br>The differential system of this model is now | |
| $\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \left[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41},t)\right]G_{40}$ | *43* |
| $\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \left[(a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41},t)\right]G_{41}$ | *44* |
| $\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \left[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41},t)\right]G_{42}$ | *45* |
| $\frac{dT_{40}}{dt} = (b_{40})^{(8)}T_{41} - \left[(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}((G_{43}),t)\right]T_{40}$ | *46* |

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| $\frac{dT_{41}}{dt} = (b_{41})^{(8)}T_{40} - \left[(b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}((G_{43}),t)\right]T_{41}$ | *47* |
| $\frac{dT_{42}}{dt} = (b_{42})^{(8)}T_{41} - \left[(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}((G_{43}),t)\right]T_{42}$ | *48* |
| **Module Numbered Nine**<br><br>The differential system of this model is now | |
| $\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \left[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45},t)\right]G_{44}$ | 49 |
| $\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \left[(a_{45}^{'})^{(9)} + (a_{45}^{''})^{(9)}(T_{45},t)\right]G_{45}$ | 50 |
| $\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \left[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45},t)\right]G_{46}$ | 51 |
| $\frac{dT_{44}}{dt} = (b_{44})^{(9)}T_{45} - \left[(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}((G_{47}),t)\right]T_{44}$ | 52 |
| $\frac{dT_{45}}{dt} = (b_{45})^{(9)}T_{44} - \left[(b_{45}^{'})^{(9)} - (b_{45}^{''})^{(9)}((G_{47}),t)\right]T_{45}$ | 53 |
| $\frac{dT_{46}}{dt} = (b_{46})^{(9)}T_{45} - \left[(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}((G_{47}),t)\right]T_{46}$ | 54 |
| $+(a_{44}^{''})^{(9)}(T_{45},t) =$ **First augmentation factor** | |
| $-(b_{44}^{''})^{(9)}((G_{47}),t) =$ **First detrition factor** | |
| $\frac{dG_{13}}{dt} = (a_{13})^{(1)}G_{14} - \begin{bmatrix} (a_{13}^{'})^{(1)}\boxed{+(a_{13}^{''})^{(1)}(T_{14},t)}\boxed{+(a_{16}^{''})^{(2,2,)}(T_{17},t)}\boxed{+(a_{20}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{24}^{''})^{(4,4,4,4,)}(T_{25},t)}\boxed{+(a_{28}^{''})^{(5,5,5,5,)}(T_{29},t)}\boxed{+(a_{32}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{36}^{''})^{(7,7)}(T_{37},t)}\boxed{+(a_{40}^{''})^{(8,8)}(T_{41},t)}\boxed{+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{13}$ | *55* |
| $\frac{dG_{14}}{dt} = (a_{14})^{(1)}G_{13} - \begin{bmatrix} (a_{14}^{'})^{(1)}\boxed{+(a_{14}^{''})^{(1)}(T_{14},t)}\boxed{+(a_{17}^{''})^{(2,2,)}(T_{17},t)}\boxed{+(a_{21}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{25}^{''})^{(4,4,4,4,)}(T_{25},t)}\boxed{+(a_{29}^{''})^{(5,5,5,5,)}(T_{29},t)}\boxed{+(a_{33}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{37}^{''})^{(7,7)}(T_{37},t)}\boxed{+(a_{41}^{''})^{(8,8)}(T_{41},t)}\boxed{+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{14}$ | *56* |
| $\frac{dG_{15}}{dt} = (a_{15})^{(1)}G_{14} - \begin{bmatrix} (a_{15}^{'})^{(1)}\boxed{+(a_{15}^{''})^{(1)}(T_{14},t)}\boxed{+(a_{18}^{''})^{(2,2,)}(T_{17},t)}\boxed{+(a_{22}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{26}^{''})^{(4,4,4,4,)}(T_{25},t)}\boxed{+(a_{30}^{''})^{(5,5,5,5,)}(T_{29},t)}\boxed{+(a_{34}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{38}^{''})^{(7,7)}(T_{37},t)}\boxed{+(a_{42}^{''})^{(8,8)}(T_{41},t)}\boxed{+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{15}$ | *57* |
| Where $\boxed{(a_{13}^{''})^{(1)}(T_{14},t)}$, $\boxed{(a_{14}^{''})^{(1)}(T_{14},t)}$, $\boxed{(a_{15}^{''})^{(1)}(T_{14},t)}$ are first augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{16}^{''})^{(2,2,)}(T_{17},t)}$, $\boxed{+(a_{17}^{''})^{(2,2,)}(T_{17},t)}$, $\boxed{+(a_{18}^{''})^{(2,2,)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{20}^{''})^{(3,3,)}(T_{21},t)}$, $\boxed{+(a_{21}^{''})^{(3,3,)}(T_{21},t)}$, $\boxed{+(a_{22}^{''})^{(3,3,)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{24}^{''})^{(4,4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{25}^{''})^{(4,4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{26}^{''})^{(4,4,4,4,)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{28}^{''})^{(5,5,5,5,)}(T_{29},t)}$, $\boxed{+(a_{29}^{''})^{(5,5,5,5,)}(T_{29},t)}$, $\boxed{+(a_{30}^{''})^{(5,5,5,5,)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{32}^{''})^{(6,6,6,6,)}(T_{33},t)}$, $\boxed{+(a_{33}^{''})^{(6,6,6,6,)}(T_{33},t)}$, $\boxed{+(a_{34}^{''})^{(6,6,6,6,)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{38}^{''})^{(7,7)}(T_{37},t)}$ $\boxed{+(a_{37}^{''})^{(7,7)}(T_{37},t)}$ $\boxed{+(a_{36}^{''})^{(7,7)}(T_{37},t)}$ are seventh augmentation coefficient for 1,2,3 | |

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| $+(a_{40}^{''})^{(8,8)}(T_{41},t)$ $+(a_{41}^{''})^{(8,8)}(T_{41},t)$ $+(a_{42}^{''})^{(8,8)}(T_{41},t)$ are eight augmentation coefficient for 1,2,3 $+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficient for 1,2,3 | |
| $\frac{dT_{13}}{dt} = (b_{13})^{(1)}T_{14} - \begin{bmatrix} (b_{13}^{'})^{(1)} & -(b_{13}^{''})^{(1)}(G,t) & -(b_{16}^{''})^{(2,2,)}(G_{19},t) & -(b_{20}^{''})^{(3,3,)}(G_{23},t) \\ -(b_{24}^{''})^{(4,4,4,4,)}(G_{27},t) & -(b_{28}^{''})^{(5,5,5,5,)}(G_{31},t) & -(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t) \\ -(b_{36}^{''})^{(7,7,)}(G_{39},t) & -(b_{40}^{''})^{(8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{13}$ | 58 |
| $\frac{dT_{14}}{dt} = (b_{14})^{(1)}T_{13} - \begin{bmatrix} (b_{14}^{'})^{(1)} & -(b_{14}^{''})^{(1)}(G,t) & -(b_{17}^{''})^{(2,2,)}(G_{19},t) & -(b_{21}^{''})^{(3,3,)}(G_{23},t) \\ -(b_{25}^{''})^{(4,4,4,4,)}(G_{27},t) & -(b_{29}^{''})^{(5,5,5,5,)}(G_{31},t) & -(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t) \\ -(b_{37}^{''})^{(7,7,)}(G_{39},t) & -(b_{41}^{''})^{(8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{14}$ | 59 |
| $\frac{dT_{15}}{dt} = (b_{15})^{(1)}T_{14} - \begin{bmatrix} (b_{15}^{'})^{(1)} & -(b_{15}^{''})^{(1)}(G,t) & -(b_{18}^{''})^{(2,2,)}(G_{19},t) & -(b_{22}^{''})^{(3,3,)}(G_{23},t) \\ -(b_{26}^{''})^{(4,4,4,4,)}(G_{27},t) & -(b_{30}^{''})^{(5,5,5,5,)}(G_{31},t) & -(b_{34}^{''})^{(6,6,6,6,)}(G_{35},t) \\ -(b_{38}^{''})^{(7,7,)}(G_{39},t) & -(b_{42}^{''})^{(8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{15}$ | *60* |
| Where $-(b_{13}^{''})^{(1)}(G,t)$, $-(b_{14}^{''})^{(1)}(G,t)$, $-(b_{15}^{''})^{(1)}(G,t)$ are first detrition coefficients for category 1, 2 and 3 $-(b_{16}^{''})^{(2,2,)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3 $-(b_{20}^{''})^{(3,3,)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3 $-(b_{24}^{''})^{(4,4,4,4,)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3 $-(b_{28}^{''})^{(5,5,5,5,)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3 $-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,)}(G_{35},t)$ are sixth detrition coefficients for category 1, 2 and 3 $-(b_{37}^{''})^{(7,7,)}(G_{39},t)$, $-(b_{36}^{''})^{(7,7,)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3 $-(b_{40}^{''})^{(8,8)}(G_{43},t)$ $-(b_{41}^{''})^{(8,8)}(G_{43},t)$ $-(b_{42}^{''})^{(8,8)}(G_{43},t)$ are eight detrition coefficients for category 1, 2 and 3 $-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3 | |
| | |
| $\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \begin{bmatrix} (a_{16}^{'})^{(2)} & +(a_{16}^{''})^{(2)}(T_{17},t) & +(a_{13}^{''})^{(1,1,)}(T_{14},t) & +(a_{20}^{''})^{(3,3,3)}(T_{21},t) \\ +(a_{24}^{''})^{(4,4,4,4,4)}(T_{25},t) & +(a_{28}^{''})^{(5,5,5,5,5)}(T_{29},t) & +(a_{32}^{''})^{(6,6,6,6,6)}(T_{33},t) \\ +(a_{36}^{''})^{(7,7,7)}(T_{37},t) & +(a_{40}^{''})^{(8,8,8)}(T_{41},t) & +(a_{44}^{''})^{(9,9)}(T_{45},t) \end{bmatrix} G_{16}$ | *61* |

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| $\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \begin{bmatrix} (a_{17}^{'})^{(2)}\boxed{+(a_{17}^{''})^{(2)}(T_{17},t)} & \boxed{+(a_{14}^{''})^{(1,1,)}(T_{14},t)} & \boxed{+(a_{21}^{''})^{(3,3,3)}(T_{21},t)} \\ \boxed{+(a_{25}^{''})^{(4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{29}^{''})^{(5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{33}^{''})^{(6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{37}^{''})^{(7,7,7)}(T_{37},t)} & \boxed{+(a_{41}^{''})^{(8,8,8)}(T_{41},t)} & \boxed{+(a_{45}^{''})^{(9,9)}(T_{45},t)} \end{bmatrix} G_{17}$ | 62 |
| $\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \begin{bmatrix} (a_{18}^{'})^{(2)}\boxed{+(a_{18}^{''})^{(2)}(T_{17},t)} & \boxed{+(a_{15}^{''})^{(1,1,)}(T_{14},t)} & \boxed{+(a_{22}^{''})^{(3,3,3)}(T_{21},t)} \\ \boxed{+(a_{26}^{''})^{(4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{30}^{''})^{(5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{34}^{''})^{(6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{38}^{''})^{(7,7,7)}(T_{37},t)} & \boxed{+(a_{42}^{''})^{(8,8,8)}(T_{41},t)} & \boxed{+(a_{46}^{''})^{(9,9)}(T_{45},t)} \end{bmatrix} G_{18}$ | 63 |
| Where $\boxed{+(a_{16}^{''})^{(2)}(T_{17},t)}$, $\boxed{+(a_{17}^{''})^{(2)}(T_{17},t)}$, $\boxed{+(a_{18}^{''})^{(2)}(T_{17},t)}$ are first augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{13}^{''})^{(1,1,)}(T_{14},t)}$, $\boxed{+(a_{14}^{''})^{(1,1,)}(T_{14},t)}$, $\boxed{+(a_{15}^{''})^{(1,1,)}(T_{14},t)}$ are second augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{20}^{''})^{(3,3,3)}(T_{21},t)}$, $\boxed{+(a_{21}^{''})^{(3,3,3)}(T_{21},t)}$, $\boxed{+(a_{22}^{''})^{(3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{24}^{''})^{(4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{25}^{''})^{(4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{26}^{''})^{(4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{28}^{''})^{(5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{29}^{''})^{(5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{30}^{''})^{(5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{32}^{''})^{(6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{33}^{''})^{(6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{34}^{''})^{(6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{36}^{''})^{(7,7,7)}(T_{37},t)}$, $\boxed{+(a_{37}^{''})^{(7,7,7)}(T_{37},t)}$, $\boxed{+(a_{38}^{''})^{(7,7,7)}(T_{37},t)}$ are seventh augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{40}^{''})^{(8,8,8)}(T_{41},t)}$, $\boxed{+(a_{41}^{''})^{(8,8,8)}(T_{41},t)}$, $\boxed{+(a_{42}^{''})^{(8,8,8)}(T_{41},t)}$ are eight augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{44}^{''})^{(9,9)}(T_{45},t)}$, $\boxed{+(a_{45}^{''})^{(9,9)}(T_{45},t)}$, $\boxed{+(a_{46}^{''})^{(9,9)}(T_{45},t)}$ are ninth augmentation coefficient for category 1, 2 and 3 | |
| $\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \begin{bmatrix} (b_{16}^{'})^{(2)}\boxed{-(b_{16}^{''})^{(2)}(G_{19},t)} & \boxed{-(b_{13}^{''})^{(1,1,)}(G,t)} & \boxed{-(b_{20}^{''})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b_{24}^{''})^{(4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{28}^{''})^{(5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{32}^{''})^{(6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{36}^{''})^{(7,7,7)}(G_{39},t)} & \boxed{-(b_{40}^{''})^{(8,8,8)}(G_{43},t)} & \boxed{-(b_{44}^{''})^{(9,9)}(G_{47},t)} \end{bmatrix} T_{16}$ | 64 |
| $\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \begin{bmatrix} (b_{17}^{'})^{(2)}\boxed{-(b_{17}^{''})^{(2)}(G_{19},t)} & \boxed{-(b_{14}^{''})^{(1,1,)}(G,t)} & \boxed{-(b_{21}^{''})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b_{25}^{''})^{(4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{29}^{''})^{(5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{33}^{''})^{(6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{37}^{''})^{(7,7,7)}(G_{39},t)} & \boxed{-(b_{41}^{''})^{(8,8,8)}(G_{43},t)} & \boxed{-(b_{45}^{''})^{(9,9)}(G_{47},t)} \end{bmatrix} T_{17}$ | 65 |
| $\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \begin{bmatrix} (b_{18}^{'})^{(2)}\boxed{-(b_{18}^{''})^{(2)}(G_{19},t)} & \boxed{-(b_{15}^{''})^{(1,1,)}(G,t)} & \boxed{-(b_{22}^{''})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b_{26}^{''})^{(4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{30}^{''})^{(5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{34}^{''})^{(6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{38}^{''})^{(7,7,7)}(G_{39},t)} & \boxed{-(b_{42}^{''})^{(8,8,8)}(G_{43},t)} & \boxed{-(b_{46}^{''})^{(9,9)}(G_{47},t)} \end{bmatrix} T_{18}$ | 66 |
| where $\boxed{-(b_{16}^{''})^{(2)}(G_{19},t)}$, $\boxed{-(b_{17}^{''})^{(2)}(G_{19},t)}$, $\boxed{-(b_{18}^{''})^{(2)}(G_{19},t)}$ are first detrition coefficients for category 1, 2 and 3 | |

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| $-(b_{13}^{''})^{(1,1,)}(G,t)$, $-(b_{14}^{''})^{(1,1,)}(G,t)$, $-(b_{15}^{''})^{(1,1,)}(G,t)$ are second detrition coefficients for category 1,2 and 3<br>$-(b_{20}^{''})^{(3,3,3,)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,)}(G_{23},t)$ are third detrition coefficients for category 1,2 and 3<br>$-(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1,2 and 3<br>$-(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1,2 and 3<br>$-(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$ are sixth detrition coefficients for category 1,2 and 3<br>$-(b_{36}^{''})^{(7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1,2 and 3<br>$-(b_{40}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8)}(G_{43},t)$ are eight detrition coefficients for category 1,2 and 3<br>$-(b_{44}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1,2 and 3 | |
| $\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \left[ \begin{array}{c} (a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21},t) + (a_{16}^{''})^{(2,2,2)}(T_{17},t) + (a_{13}^{''})^{(1,1,1,)}(T_{14},t) \\ + (a_{24}^{''})^{(4,4,4,4,4,4)}(T_{25},t) + (a_{28}^{''})^{(5,5,5,5,5,5)}(T_{29},t) + (a_{32}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ + (a_{36}^{''})^{(7,7,7,7)}(T_{37},t) + (a_{40}^{''})^{(8,8,8,8)}(T_{41},t) + (a_{44}^{''})^{(9,9,9)}(T_{45},t) \end{array} \right] G_{20}$ | 67 |
| $\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \left[ \begin{array}{c} (a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)}(T_{21},t) + (a_{17}^{''})^{(2,2,2)}(T_{17},t) + (a_{14}^{''})^{(1,1,1,)}(T_{14},t) \\ + (a_{25}^{''})^{(4,4,4,4,4,4)}(T_{25},t) + (a_{29}^{''})^{(5,5,5,5,5,5)}(T_{29},t) + (a_{33}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ + (a_{37}^{''})^{(7,7,7,7)}(T_{37},t) + (a_{41}^{''})^{(8,8,8,8)}(T_{41},t) + (a_{45}^{''})^{(9,9,9)}(T_{45},t) \end{array} \right] G_{21}$ | 68 |
| $\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \left[ \begin{array}{c} (a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21},t) + (a_{18}^{''})^{(2,2,2)}(T_{17},t) + (a_{15}^{''})^{(1,1,1,)}(T_{14},t) \\ + (a_{26}^{''})^{(4,4,4,4,4,4)}(T_{25},t) + (a_{30}^{''})^{(5,5,5,5,5,5)}(T_{29},t) + (a_{34}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ + (a_{38}^{''})^{(7,7,7,7)}(T_{37},t) + (a_{42}^{''})^{(8,8,8,8)}(T_{41},t) + (a_{46}^{''})^{(9,9,9)}(T_{45},t) \end{array} \right] G_{22}$ | 69 |
| $+(a_{20}^{''})^{(3)}(T_{21},t)$, $+(a_{21}^{''})^{(3)}(T_{21},t)$, $+(a_{22}^{''})^{(3)}(T_{21},t)$ are first augmentation coefficients for category 1, 2 and 3<br>$+(a_{16}^{''})^{(2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2)}(T_{17},t)$ are second augmentation coefficients for category 1, 2 and 3<br>$+(a_{13}^{''})^{(1,1,1,)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,)}(T_{14},t)$ are third augmentation coefficients for category 1, 2 and 3<br>$+(a_{24}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficients for category 1, 2 and 3<br>$+(a_{28}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficients for category 1, 2 and 3<br>$+(a_{32}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation | |

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| coefficients for category 1, 2 and 3 $+(a_{36}^{''})^{(7,7,7,7)}(T_{37},t)$, $+(a_{37}^{''})^{(7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7)}(T_{37},t)$ are seventh augmentation coefficients for category 1, 2 and 3 $+(a_{40}^{''})^{(8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8)}(T_{41},t)$ are eight augmentation coefficients for category 1, 2 and 3 $+(a_{44}^{''})^{(9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9)}(T_{45},t)$, $+(a_{46}^{''})^{(9,9,9)}(T_{45},t)$ are ninth augmentation coefficients for category 1, 2 and 3 | |
| $\frac{dT_{20}}{dt}=(b_{20})^{(3)}T_{21}-\begin{bmatrix}(b_{20}^{'})^{(3)}-(b_{20}^{''})^{(3)}(G_{23},t) & -(b_{16}^{''})^{(2,2,2)}(G_{19},t) & -(b_{13}^{''})^{(1,1,1,)}(G,t)\\ -(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t) & -(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t) & -(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t)\\ -(b_{36}^{''})^{(7,7,7,7)}(G_{39},t) & -(b_{40}^{''})^{(8,8,8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9,9)}(G_{47},t)\end{bmatrix}T_{20}$ | *70* |
| $\frac{dT_{21}}{dt}=(b_{21})^{(3)}T_{20}-\begin{bmatrix}(b_{21}^{'})^{(3)}-(b_{21}^{''})^{(3)}(G_{23},t) & -(b_{17}^{''})^{(2,2,2)}(G_{19},t) & -(b_{14}^{''})^{(1,1,1,)}(G,t)\\ -(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t) & -(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t) & -(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t)\\ -(b_{37}^{''})^{(7,7,7,7)}(G_{39},t) & -(b_{41}^{''})^{(8,8,8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9,9)}(G_{47},t)\end{bmatrix}T_{21}$ | *71* |
| $\frac{dT_{22}}{dt}=(b_{22})^{(3)}T_{21}-\begin{bmatrix}(b_{22}^{'})^{(3)}-(b_{22}^{''})^{(3)}(G_{23},t) & -(b_{18}^{''})^{(2,2,2)}(G_{19},t) & -(b_{15}^{''})^{(1,1,1,)}(G,t)\\ -(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t) & -(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t) & -(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t)\\ -(b_{38}^{''})^{(7,7,7,7)}(G_{39},t) & -(b_{42}^{''})^{(8,8,8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9)}(G_{47},t)\end{bmatrix}T_{22}$ | *72* |
| $-(b_{20}^{''})^{(3)}(G_{23},t)$, $-(b_{21}^{''})^{(3)}(G_{23},t)$, $-(b_{22}^{''})^{(3)}(G_{23},t)$ *are first detrition coefficients for category 1, 2 and 3* $-(b_{16}^{''})^{(2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2)}(G_{19},t)$ *are second detrition coefficients for category 1, 2 and 3* $-(b_{13}^{''})^{(1,1,1,)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,)}(G,t)$, $-(b_{15}^{''})^{(1,1,1,)}(G,t)$ *are third detrition coefficients for category 1,2 and 3* $-(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$ *are fourth detrition coefficients for category 1, 2 and 3* $-(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$ *are fifth detrition coefficients for category 1, 2 and 3* $-(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$ *are sixth detrition coefficients for category 1, 2 and 3* $-(b_{36}^{''})^{(7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3 $-(b_{40}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8)}(G_{43},t)$ are eight detrition coefficients for category 1, 2 and 3 $-(b_{46}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3 | |

$$\frac{dG_{24}}{dt} = (a_{24})^{(4)}G_{25} - \begin{bmatrix} (a'_{24})^{(4)} & +(a''_{24})^{(4)}(T_{25},t) & +(a''_{28})^{(5,5,)}(T_{29},t) & +(a''_{32})^{(6,6,)}(T_{33},t) \\ & +(a''_{13})^{(1,1,1,1)}(T_{14},t) & +(a''_{16})^{(2,2,2,2)}(T_{17},t) & +(a''_{20})^{(3,3,3,3)}(T_{21},t) \\ & +(a''_{36})^{(7,7,7,7,7)}(T_{37},t) & +(a''_{40})^{(8,8,8,8,8)}(T_{41},t) & +(a''_{44})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{24} \quad (73)$$

$$\frac{dG_{25}}{dt} = (a_{25})^{(4)}G_{24} - \begin{bmatrix} (a'_{25})^{(4)} & +(a''_{25})^{(4)}(T_{25},t) & +(a''_{29})^{(5,5,)}(T_{29},t) & +(a''_{33})^{(6,6)}(T_{33},t) \\ & +(a''_{14})^{(1,1,1,1)}(T_{14},t) & +(a''_{17})^{(2,2,2,2)}(T_{17},t) & +(a''_{21})^{(3,3,3,3)}(T_{21},t) \\ & +(a''_{37})^{(7,7,7,7,7)}(T_{37},t) & +(a''_{41})^{(8,8,8,8,8)}(T_{41},t) & +(a''_{45})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{25} \quad (74)$$

$$\frac{dG_{26}}{dt} = (a_{26})^{(4)}G_{25} - \begin{bmatrix} (a'_{26})^{(4)} & +(a''_{26})^{(4)}(T_{25},t) & +(a''_{30})^{(5,5,)}(T_{29},t) & +(a''_{34})^{(6,6,)}(T_{33},t) \\ & +(a''_{15})^{(1,1,1,1)}(T_{14},t) & +(a''_{18})^{(2,2,2,2)}(T_{17},t) & +(a''_{22})^{(3,3,3,3)}(T_{21},t) \\ & +(a''_{38})^{(7,7,7,7,7)}(T_{37},t) & +(a''_{42})^{(8,8,8,8,8)}(T_{41},t) & +(a''_{46})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{26} \quad (75)$$

$(a''_{24})^{(4)}(T_{25},t)$, $(a''_{25})^{(4)}(T_{25},t)$, $(a''_{26})^{(4)}(T_{25},t)$ are first augmentation coefficients category 1, 2 3

$+(a''_{28})^{(5,5,)}(T_{29},t)$, $+(a''_{29})^{(5,5,)}(T_{29},t)$, $+(a''_{30})^{(5,5,)}(T_{29},t)$ are second augmentation coefficient for category 1, 2 and 3

$+(a''_{32})^{(6,6,)}(T_{33},t)$, $+(a''_{33})^{(6,6,)}(T_{33},t)$, $+(a''_{34})^{(6,6,)}(T_{33},t)$ are third augmentation coefficient for category 1, 2 and 3

$+(a''_{13})^{(1,1,1,1)}(T_{14},t)$, $+(a''_{14})^{(1,1,1,1)}(T_{14},t)$, $+(a''_{15})^{(1,1,1,1)}(T_{14},t)$ are fourth augmentation coefficients for category 1, 2 and 3

$+(a''_{16})^{(2,2,2,2)}(T_{17},t)$, $+(a''_{17})^{(2,2,2,2)}(T_{17},t)$, $+(a''_{18})^{(2,2,2,2)}(T_{17},t)$ are fifth augmentation coefficients for category 1, 2 and 3

$+(a''_{20})^{(3,3,3,3)}(T_{21},t)$, $+(a''_{21})^{(3,3,3,3)}(T_{21},t)$, $+(a''_{22})^{(3,3,3,3)}(T_{21},t)$ are sixth augmentation coefficients for category 1, 2 and 3

$+(a''_{36})^{(7,7,7,7,7)}(T_{37},t)$, $+(a''_{37})^{(7,7,7,7,7)}(T_{37},t)$, $+(a''_{38})^{(7,7,7,7,7)}(T_{37},t)$ are seventh augmentation coefficients for category 1, 2 and 3

$+(a''_{40})^{(8,8,8,8,8)}(T_{41},t)$, $+(a''_{41})^{(8,8,8,8,8)}(T_{41},t)$, $+(a''_{42})^{(8,8,8,8,8)}(T_{41},t)$ are eighth augmentation coefficients for category 1, 2 and 3

$+(a''_{46})^{(9,9,9,9)}(T_{45},t)$, $+(a''_{45})^{(9,9,9,9)}(T_{45},t)$, $+(a''_{44})^{(9,9,9,9)}(T_{45},t)$ are ninth detrition coefficients for category 1 2 3

$$\frac{dT_{24}}{dt} = (b_{24})^{(4)}T_{25} - \begin{bmatrix} (b'_{24})^{(4)} & -(b''_{24})^{(4)}(G_{27},t) & -(b''_{28})^{(5,5,)}(G_{31},t) & -(b''_{32})^{(6,6,)}(G_{35},t) \\ & -(b''_{13})^{(1,1,1,1)}(G,t) & -(b''_{16})^{(2,2,2,2)}(G_{19},t) & -(b''_{20})^{(3,3,3,3)}(G_{23},t) \\ & -(b''_{36})^{(7,7,7,7,7)}(G_{39},t) & -(b''_{40})^{(8,8,8,8,8)}(G_{43},t) & -(b''_{44})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{24} \quad (76)$$

$$\frac{dT_{25}}{dt} = (b_{25})^{(4)}T_{24} - \begin{bmatrix} (b'_{25})^{(4)} & -(b''_{25})^{(4)}(G_{27},t) & -(b''_{29})^{(5,5,)}(G_{31},t) & -(b''_{33})^{(6,6,)}(G_{35},t) \\ & -(b''_{14})^{(1,1,1,1)}(G,t) & -(b''_{17})^{(2,2,2,2)}(G_{19},t) & -(b''_{21})^{(3,3,3,3)}(G_{23},t) \\ & -(b''_{37})^{(7,7,7,7,7)}(G_{39},t) & -(b''_{41})^{(8,8,8,8,8)}(G_{43},t) & -(b''_{45})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{25} \quad (77)$$

$$\frac{dT_{26}}{dt} = (b_{26})^{(4)}T_{25} - \begin{bmatrix} (b'_{26})^{(4)}\boxed{-(b''_{26})^{(4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,)}(G_{35},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1)}(G,t)} & \boxed{-(b''_{18})^{(2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{38})^{(7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{26} \quad (78)$$

Where $\boxed{-(b''_{24})^{(4)}(G_{27},t)}$, $\boxed{-(b''_{25})^{(4)}(G_{27},t)}$, $\boxed{-(b''_{26})^{(4)}(G_{27},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{28})^{(5,5,)}(G_{31},t)}$, $\boxed{-(b''_{29})^{(5,5,)}(G_{31},t)}$, $\boxed{-(b''_{30})^{(5,5,)}(G_{31},t)}$ are second detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{32})^{(6,6,)}(G_{35},t)}$, $\boxed{-(b''_{33})^{(6,6,)}(G_{35},t)}$, $\boxed{-(b''_{34})^{(6,6,)}(G_{35},t)}$ are third detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{13})^{(1,1,1,1)}(G,t)}$, $\boxed{-(b''_{14})^{(1,1,1,1)}(G,t)}$, $\boxed{-(b''_{15})^{(1,1,1,1)}(G,t)}$ are fourth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{16})^{(2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{17})^{(2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{18})^{(2,2,2,2)}(G_{19},t)}$ are fifth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{20})^{(3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{21})^{(3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{22})^{(3,3,3,3)}(G_{23},t)}$ are sixth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{36})^{(7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{37})^{(7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{38})^{(7,7,7,7,7)}(G_{39},t)}$ are seventh detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{40})^{(8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{41})^{(8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{42})^{(8,8,8,8,8)}(G_{43},t)}$ are eighth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{46})^{(9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{45})^{(9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{44})^{(9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1 2 3

$$\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \begin{bmatrix} (a'_{28})^{(5)}\boxed{+(a''_{28})^{(5)}(T_{29},t)} & \boxed{+(a''_{24})^{(4,4,)}(T_{25},t)} & \boxed{+(a''_{32})^{(6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{36})^{(7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{40})^{(8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{44})^{(9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{28} \quad (79)$$

$$\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \begin{bmatrix} (a'_{29})^{(5)}\boxed{+(a''_{29})^{(5)}(T_{29},t)} & \boxed{+(a''_{25})^{(4,4,)}(T_{25},t)} & \boxed{+(a''_{33})^{(6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{14})^{(1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{17})^{(2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{21})^{(3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{37})^{(7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{41})^{(8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{45})^{(9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{29} \quad (80)$$

$$\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \begin{bmatrix} (a'_{30})^{(5)}\boxed{+(a''_{30})^{(5)}(T_{29},t)} & \boxed{+(a''_{26})^{(4,4,)}(T_{25},t)} & \boxed{+(a''_{34})^{(6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{15})^{(1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{18})^{(2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{22})^{(3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{38})^{(7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{42})^{(8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{46})^{(9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{30} \quad (81)$$

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| *Where* $+(a_{28}^{''})^{(5)}(T_{29},t)$ , $+(a_{29}^{''})^{(5)}(T_{29},t)$ , $+(a_{30}^{''})^{(5)}(T_{29},t)$ *are first augmentation coefficients for category* 1, 2 *and* 3<br>*And* $+(a_{24}^{''})^{(4,4,)}(T_{25},t)$ , $+(a_{25}^{''})^{(4,4,)}(T_{25},t)$ , $+(a_{26}^{''})^{(4,4,)}(T_{25},t)$ *are second augmentation coefficient for category* 1, 2 *and* 3<br>$+(a_{32}^{''})^{(6,6,6)}(T_{33},t)$ , $+(a_{33}^{''})^{(6,6,6)}(T_{33},t)$ , $+(a_{34}^{''})^{(6,6,6)}(T_{33},t)$ *are third augmentation coefficient for category* 1, 2 *and* 3<br>$+(a_{13}^{''})^{(1,1,1,1,1)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1,1)}(T_{14},t)$ are fourth augmentation coefficients for category 1,2, and 3<br>$+(a_{16}^{''})^{(2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2)}(T_{17},t)$ are fifth augmentation coefficients for category 1,2,and 3<br>$+(a_{20}^{''})^{(3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3)}(T_{21},t)$ are sixth augmentation coefficients for category 1,2, 3<br>$+(a_{36}^{''})^{(7,7,7,7,7,7)}(T_{37},t)$ , $+(a_{37}^{''})^{(7,7,7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7,7,7)}(T_{37},t)$ are seventh augmentation coefficients for category 1,2, 3<br>$+(a_{40}^{''})^{(8,8\ ,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8,8,8)}(T_{41},t)$ are eighth augmentation coefficients for category 1,2, 3<br>$+(a_{46}^{''})^{(9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficients for category 1,2, 3 | |
| $\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \begin{bmatrix} (b_{28}^{'})^{(5)}\boxed{-(b_{28}^{''})^{(5)}(G_{31},t)} & \boxed{-(b_{24}^{''})^{(4,4,)}(G_{27},t)} & \boxed{-(b_{32}^{''})^{(6,6,6)}(G_{35},t)} \\ \boxed{-(b_{13}^{''})^{(1,1,1,1,1)}(G,t)} & \boxed{-(b_{16}^{''})^{(2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{20}^{''})^{(3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{36}^{''})^{(7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{40}^{''})^{(8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b_{44}^{''})^{(9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{28}$ | *82* |
| $\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \begin{bmatrix} (b_{29}^{'})^{(5)}\boxed{-(b_{29}^{''})^{(5)}(G_{31},t)} & \boxed{-(b_{25}^{''})^{(4,4,)}(G_{27},t)} & \boxed{-(b_{33}^{''})^{(6,6,6)}(G_{35},t)} \\ \boxed{-(b_{14}^{''})^{(1,1,1,1,1)}(G,t)} & \boxed{-(b_{17}^{''})^{(2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{21}^{''})^{(3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{37}^{''})^{(7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{41}^{''})^{(8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b_{45}^{''})^{(9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{29}$ | *83* |
| $\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \begin{bmatrix} (b_{30}^{'})^{(5)}\boxed{-(b_{30}^{''})^{(5)}(G_{31},t)} & \boxed{-(b_{26}^{''})^{(4,4,)}(G_{27},t)} & \boxed{-(b_{34}^{''})^{(6,6,6)}(G_{35},t)} \\ \boxed{-(b_{15}^{''})^{(1,1,1,1,1,)}(G,t)} & \boxed{-(b_{18}^{''})^{(2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{22}^{''})^{(3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{38}^{''})^{(7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{42}^{''})^{(8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b_{46}^{''})^{(9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{30}$ | *84* |
| *where* $\boxed{-(b_{28}^{''})^{(5)}(G_{31},t)}$ , $\boxed{-(b_{29}^{''})^{(5)}(G_{31},t)}$ , $\boxed{-(b_{30}^{''})^{(5)}(G_{31},t)}$ *are first detrition coefficients for category* 1, 2 *and* 3<br>$\boxed{-(b_{24}^{''})^{(4,4,)}(G_{27},t)}$, $\boxed{-(b_{25}^{''})^{(4,4,)}(G_{27},t)}$, $\boxed{-(b_{26}^{''})^{(4,4,)}(G_{27},t)}$ *are second detrition coefficients for category* 1,2 *and* 3<br>$\boxed{-(b_{32}^{''})^{(6,6,6)}(G_{35},t)}$, $\boxed{-(b_{33}^{''})^{(6,6,6)}(G_{35},t)}$, $\boxed{-(b_{34}^{''})^{(6,6,6)}(G_{35},t)}$ *are third detrition coefficients for category* 1,2 *and* 3<br>$\boxed{-(b_{13}^{''})^{(1,1,1,1,1)}(G,t)}$, $\boxed{-(b_{14}^{''})^{(1,1,1,1,1)}(G,t)}$ , $\boxed{-(b_{15}^{''})^{(1,1,1,1,1,)}(G,t)}$ are fourth detrition coefficients for category 1,2, and 3<br>$\boxed{-(b_{16}^{''})^{(2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b_{17}^{''})^{(2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b_{18}^{''})^{(2,2,2,2,2)}(G_{19},t)}$ are fifth detrition coefficients | |

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| for category 1,2, and 3<br>$\boxed{-(b_{20}^{''})^{(3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b_{21}^{''})^{(3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b_{22}^{''})^{(3,3,3,3,3)}(G_{23},t)}$ are sixth detrition coefficients for category 1,2, and 3<br>$\boxed{-(b_{36}^{''})^{(7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{37}^{''})^{(7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{38}^{''})^{(7,7,7,7,7,7)}(G_{39},t)}$ are seventh detrition coefficients for category 1,2, and 3<br>$\boxed{-(b_{42}^{''})^{(8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{41}^{''})^{(8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{40}^{''})^{(8,8,8,8,8,8)}(G_{43},t)}$ are eighth detrition coefficients for category 1,2, and 3<br>$\boxed{-(b_{46}^{''})^{(9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b_{45}^{''})^{(9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b_{44}^{''})^{(9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1,2, and 3 | |
| $\frac{dG_{32}}{dt} = (a_{32})^{(6)}G_{33} - \begin{bmatrix} (a_{32}^{'})^{(6)}\boxed{+(a_{32}^{''})^{(6)}(T_{33},t)} & \boxed{+(a_{28}^{''})^{(5,5,5)}(T_{29},t)} & \boxed{+(a_{24}^{''})^{(4,4,4,)}(T_{25},t)} \\ \boxed{+(a_{13}^{''})^{(1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{16}^{''})^{(2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{20}^{''})^{(3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{36}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{40}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{44}^{''})^{(9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{32}$ | *85* |
| $\frac{dG_{33}}{dt} = (a_{33})^{(6)}G_{32} - \begin{bmatrix} (a_{33}^{'})^{(6)}\boxed{+(a_{33}^{''})^{(6)}(T_{33},t)} & \boxed{+(a_{29}^{''})^{(5,5,5)}(T_{29},t)} & \boxed{+(a_{25}^{''})^{(4,4,4,)}(T_{25},t)} \\ \boxed{+(a_{14}^{''})^{(1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{17}^{''})^{(2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{21}^{''})^{(3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{37}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{41}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{45}^{''})^{(9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{33}$ | *86* |
| $\frac{dG_{34}}{dt} = (a_{34})^{(6)}G_{33} - \begin{bmatrix} (a_{34}^{'})^{(6)}\boxed{+(a_{34}^{''})^{(6)}(T_{33},t)} & \boxed{+(a_{30}^{''})^{(5,5,5)}(T_{29},t)} & \boxed{+(a_{26}^{''})^{(4,4,4,)}(T_{25},t)} \\ \boxed{+(a_{15}^{''})^{(1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{18}^{''})^{(2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{22}^{''})^{(3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{38}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{42}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{46}^{''})^{(9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{34}$ | *87* |
| $\boxed{+(a_{32}^{''})^{(6)}(T_{33},t)}$, $\boxed{+(a_{33}^{''})^{(6)}(T_{33},t)}$, $\boxed{+(a_{34}^{''})^{(6)}(T_{33},t)}$ *are first augmentation coefficients for category* 1, 2 *and* 3<br>$\boxed{+(a_{28}^{''})^{(5,5,5)}(T_{29},t)}$, $\boxed{+(a_{29}^{''})^{(5,5,5)}(T_{29},t)}$, $\boxed{+(a_{30}^{''})^{(5,5,5)}(T_{29},t)}$ *are second augmentation coefficients for category* 1, 2 *and* 3<br>$\boxed{+(a_{24}^{''})^{(4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{25}^{''})^{(4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{26}^{''})^{(4,4,4,)}(T_{25},t)}$ *are third augmentation coefficients for category* 1, 2 *and* 3<br>$\boxed{+(a_{13}^{''})^{(1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a_{14}^{''})^{(1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a_{15}^{''})^{(1,1,1,1,1,1)}(T_{14},t)}$ - are fourth augmentation coefficients<br>$\boxed{+(a_{16}^{''})^{(2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a_{17}^{''})^{(2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a_{18}^{''})^{(2,2,2,2,2,2)}(T_{17},t)}$ - fifth augmentation coefficients<br>$\boxed{+(a_{20}^{''})^{(3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a_{21}^{''})^{(3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a_{22}^{''})^{(3,3,3,3,3,3)}(T_{21},t)}$ sixth augmentation coefficients<br>$\boxed{+(a_{36}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)}$, $\boxed{+(a_{37}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)}$, $\boxed{+(a_{38}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)}$ seventh augmentation coefficients<br>$\boxed{+(a_{40}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{41}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{42}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)}$ Eighth augmentation coefficients<br>$\boxed{+(a_{44}^{''})^{(9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{45}^{''})^{(9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{46}^{''})^{(9,9,9,9,9,9)}(T_{45},t)}$ ninth augmentation | |

coefficients

$$\frac{dT_{32}}{dt} = (b_{32})^{(6)}T_{33} - \begin{bmatrix} (b'_{32})^{(6)}\boxed{-(b''_{32})^{(6)}(G_{35},t)} & \boxed{-(b''_{28})^{(5,5,5)}(G_{31},t)} & \boxed{-(b''_{24})^{(4,4,4,)}(G_{27},t)} \\ \boxed{-(b''_{13})^{(1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{16})^{(2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{20})^{(3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{36})^{(7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{40})^{(8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{44})^{(9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{32} \quad (88)$$

$$\frac{dT_{33}}{dt} = (b_{33})^{(6)}T_{32} - \begin{bmatrix} (b'_{33})^{(6)}\boxed{-(b''_{33})^{(6)}(G_{35},t)} & \boxed{-(b''_{29})^{(5,5,5)}(G_{31},t)} & \boxed{-(b''_{25})^{(4,4,4,)}(G_{27},t)} \\ \boxed{-(b''_{14})^{(1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{17})^{(2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{21})^{(3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{37})^{(7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{41})^{(8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{45})^{(9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{33} \quad (89)$$

$$\frac{dT_{34}}{dt} = (b_{34})^{(6)}T_{33} - \begin{bmatrix} (b'_{34})^{(6)}\boxed{-(b''_{34})^{(6)}(G_{35},t)} & \boxed{-(b''_{30})^{(5,5,5)}(G_{31},t)} & \boxed{-(b''_{26})^{(4,4,4,)}(G_{27},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{18})^{(2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{38})^{(7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{34} \quad (90)$$

$\boxed{-(b''_{32})^{(6)}(G_{35},t)}, \boxed{-(b''_{33})^{(6)}(G_{35},t)}, \boxed{-(b''_{34})^{(6)}(G_{35},t)}$ *are first detrition coefficients for category* 1, 2 *and* 3

$\boxed{-(b''_{28})^{(5,5,5)}(G_{31},t)}, \boxed{-(b''_{29})^{(5,5,5)}(G_{31},t)}, \boxed{-(b''_{30})^{(5,5,5)}(G_{31},t)}$ *are second detrition coefficients for category* 1, 2 *and* 3

$\boxed{-(b''_{24})^{(4,4,4,)}(G_{27},t)}, \boxed{-(b''_{25})^{(4,4,4,)}(G_{27},t)}, \boxed{-(b''_{26})^{(4,4,4,)}(G_{27},t)}$ *are third detrition coefficients for category* 1,2 *and* 3

$\boxed{-(b''_{13})^{(1,1,1,1,1,1)}(G,t)}, \boxed{-(b''_{14})^{(1,1,1,1,1,1)}(G,t)}, \boxed{-(b''_{15})^{(1,1,1,1,1,1)}(G,t)}$ are fourth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{16})^{(2,2,2,2,2,2)}(G_{19},t)}, \boxed{-(b''_{17})^{(2,2,2,2,2,2)}(G_{19},t)}, \boxed{-(b''_{18})^{(2,2,2,2,2,2)}(G_{19},t)}$ are fifth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{20})^{(3,3,3,3,3,3)}(G_{23},t)}, \boxed{-(b''_{21})^{(3,3,3,3,3,3)}(G_{23},t)}, \boxed{-(b''_{22})^{(3,3,3,3,3,3)}(G_{23},t)}$ are sixth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{36})^{(7,7,7,7,7,7,7)}(G_{39},t)}, \boxed{-(b''_{37})^{(7,7,7,7,7,7,7)}(G_{39},t)}, \boxed{-(b''_{38})^{(7,7,7,7,7,7,7)}(G_{39},t)}$ are seventh detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{40})^{(8,8,8,8,8,8,8)}(G_{43},t)}, \boxed{-(b''_{41})^{(8,8,8,8,8,8,8)}(G_{43},t)}, \boxed{-(b''_{42})^{(8,8,8,8,8,8,8)}(G_{43},t)}$ are eighth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{46})^{(9,9,9,9,9,9)}(G_{47},t)}, \boxed{-(b''_{45})^{(9,9,9,9,9,9)}(G_{47},t)}, \boxed{-(b''_{44})^{(9,9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2, and 3

$$\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \begin{bmatrix} (a'_{36})^{(7)}\boxed{+(a''_{36})^{(7)}(T_{37},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{24})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8,)}(T_{41},t)} & \boxed{+(a''_{44})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{13} \quad (91)$$

$$\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \begin{bmatrix} (a_{37}')^{(7)}\boxed{+(a_{37}'')^{(7)}(T_{37},t)} & \boxed{+(a_{17}'')^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{21}'')^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{25}'')^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{29}'')^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{33}'')^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{13}'')^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{41}'')^{(8,8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{45}'')^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{14} \quad (92)$$

$$\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \begin{bmatrix} (a_{38}')^{(7)}\boxed{+(a_{38}'')^{(7)}(T_{37},t)} & \boxed{+(a_{18}'')^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{22}'')^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{26}'')^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{30}'')^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{34}'')^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{15}'')^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{42}'')^{(8,8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{46}'')^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{15} \quad (93)$$

Where $\boxed{(a_{36}'')^{(7)}(T_{37},t)}$, $\boxed{(a_{37}'')^{(7)}(T_{37},t)}$, $\boxed{(a_{38}'')^{(7)}(T_{37},t)}$ are first augmentation coefficients for category 1, 2 and 3

$\boxed{+(a_{16}'')^{(2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a_{17}'')^{(2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a_{18}'')^{(2,2,2,2,2,2,2)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{20}'')^{(3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a_{21}'')^{(3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a_{22}'')^{(3,3,3,3,3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{24}'')^{(4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{25}'')^{(4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{26}'')^{(4,4,4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{28}'')^{(5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{29}'')^{(5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{30}'')^{(5,5,5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{32}'')^{(6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{33}'')^{(6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{34}'')^{(6,6,6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{13}'')^{(1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a_{13}'')^{(1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a_{15}'')^{(1,1,1,1,1,1,1)}(T_{14},t)}$ are seventh augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{42}'')^{(8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{41}'')^{(8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{40}'')^{(8,8,8,8,8,8,8,8,)}(T_{41},t)}$ are eighth augmentation coefficient for 1,2,3

$\boxed{+(a_{46}'')^{(9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{45}'')^{(9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{44}'')^{(9,9,9,9,9,9,9)}(T_{45},t)}$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{36}}{dt} = (b_{36})^{(7)}T_{37} - \begin{bmatrix} (b_{36}')^{(7)}\boxed{-(b_{36}'')^{(7)}(G_{39},t)} & \boxed{-(b_{16}'')^{(2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{20}'')^{(3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{24}'')^{(4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{28}'')^{(5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{32}'')^{(6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{13}'')^{(1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{40}'')^{(8,8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b_{44}'')^{(9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{13} \quad (94)$$

$$\frac{dT_{37}}{dt} =$$

$$(b_{37})^{(7)}T_{36} - \begin{bmatrix} (b'_{37})^{(7)}\boxed{-(b''_{37})^{(7)}(G_{39},t)} & \boxed{-(b''_{17})^{(2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{21})^{(3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{25})^{(4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{29})^{(5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{33})^{(6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{14})^{(1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{41})^{(8,8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{45})^{(9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{14}$$

$$\frac{dT_{38}}{dt} =$$

$$(b_{38})^{(7)}T_{37} - \begin{bmatrix} (b'_{38})^{(7)}\boxed{-(b''_{38})^{(7)}(G_{39},t)} & \boxed{-(b''_{18})^{(2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{26})^{(4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{42})^{(8,8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{15}$$

Where $\boxed{-(b''_{36})^{(7)}(G_{39},t)}$, $\boxed{-(b''_{37})^{(7)}(G_{39},t)}$, $\boxed{-(b''_{38})^{(7)}(G_{39},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{16})^{(2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{17})^{(2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{18})^{(2,2,2,2,2,2,2)}(G_{19},t)}$ are second detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{20})^{(3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{21})^{(3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{22})^{(3,3,3,3,3,3,3)}(G_{23},t)}$ are third detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{24})^{(4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{25})^{(4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{26})^{(4,4,4,4,4,4,4)}(G_{27},t)}$ are fourth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{28})^{(5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{29})^{(5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{30})^{(5,5,5,5,5,5,5)}(G_{31},t)}$ are fifth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{32})^{(6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b''_{33})^{(6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b''_{34})^{(6,6,6,6,6,6,6)}(G_{35},t)}$ are sixth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{15})^{(1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{14})^{(1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{13})^{(1,1,1,1,1,1,1)}(G,t)}$ are seventh detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{40})^{(8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{41})^{(8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{42})^{(8,8,8,8,8,8,8,8)}(G_{43},t)}$ are eighth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{46})^{(9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{45})^{(9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{44})^{(9,9,9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2 and 3

$$\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \begin{bmatrix} (a'_{40})^{(8)}\boxed{+(a''_{40})^{(8)}(T_{41},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{24})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{36})^{(7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{44})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{13}$$

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$$\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \begin{bmatrix} (a'_{41})^{(8)}\boxed{+(a''_{41})^{(8)}(T_{41},t)} & \boxed{+(a''_{17})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{21})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{25})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{29})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{33})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{37})^{(7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{45})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{14}$$

$$\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \begin{bmatrix} (a'_{42})^{(8)}\boxed{+(a''_{42})^{(8)}(T_{41},t)} & \boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{46})^{(9,9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{15}$$

Where $\boxed{+(a''_{40})^{(8)}(T_{41},t)}$, $\boxed{+(a''_{41})^{(8)}(T_{41},t)}$, $\boxed{+(a''_{42})^{(8)}(T_{41},t)}$ are first augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{16})^{(2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{17})^{(2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{20})^{(3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{21})^{(3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{24})^{(4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{25})^{(4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{28})^{(5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{29})^{(5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{32})^{(6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{33})^{(6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{13})^{(1,1,1,1,1,1,1,1)}(T_{14},t)}$ $\boxed{+(a''_{14})^{(1,1,1,1,1,1,1,1)}(T_{14},t)}$ $\boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1)}(T_{14},t)}$ are seventh augmentation coefficient for 1,2,3

$\boxed{+(a''_{36})^{(7,7,7,7,7,7,7,7)}(T_{37},t)}$, $\boxed{+(a''_{37})^{(7,7,7,7,7,7,7,7)}(T_{37},t)}$, $\boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7)}(T_{37},t)}$ are eighth augmentation coefficient for 1,2,3

$\boxed{+(a''_{46})^{(9,9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a''_{45})^{(9,9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a''_{44})^{(9,9,9,9,9,9,9,9)}(T_{45},t)}$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{40}}{dt} = (b_{40})^{(8)}T_{41} - \begin{bmatrix} (b'_{40})^{(8)}\boxed{-(b''_{40})^{(8)}(G_{43},t)} & \boxed{-(b''_{16})^{(2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{20})^{(3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{24})^{(4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{28})^{(5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{32})^{(6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{13})^{(1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{36})^{(7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{44})^{(9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{13}$$

$$\frac{dT_{41}}{dt} = (b_{41})^{(8)}T_{40} - \begin{bmatrix} (b'_{41})^{(8)}\boxed{-(b''_{41})^{(8)}(G_{43},t)} & \boxed{-(b''_{17})^{(2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{21})^{(3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{25})^{(4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{29})^{(5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{33})^{(6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{14})^{(1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{37})^{(7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{45})^{(9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{14}$$

$$\frac{dT_{42}}{dt} =$$

$$(b_{42})^{(8)}T_{41} - \begin{bmatrix} (b'_{42})^{(8)}\boxed{-(b''_{42})^{(8)}(G_{43},t)} & \boxed{-(b''_{18})^{(2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{26})^{(4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{38})^{(7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{46})^{(9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{15}$$

Where $\boxed{-(b''_{36})^{(7)}(G_{39},t)}$, $\boxed{-(b''_{37})^{(7)}(G_{39},t)}$, $\boxed{-(b''_{38})^{(7)}(G_{39},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{16})^{(2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{17})^{(2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{18})^{(2,2,2,2,2,2,2,2)}(G_{19},t)}$ are second detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{20})^{(3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{21})^{(3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{22})^{(3,3,3,3,3,3,3,3)}(G_{23},t)}$ are third detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{24})^{(4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{25})^{(4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{26})^{(4,4,4,4,4,4,4,4)}(G_{27},t)}$ are fourth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{28})^{(5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{29})^{(5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{30})^{(5,5,5,5,5,5,5,5)}(G_{31},t)}$ are fifth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{32})^{(6,6,6,6,)}(G_{35},t)}$, $\boxed{-(b''_{33})^{(6,6,6,6,)}(G_{35},t)}$, $\boxed{-(b''_{15})^{(1,1,1,1,1,1,1,1)}(G,t)}$ are sixth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{13})^{(1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{14})^{(1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{38})^{(7,7,)}(G_{39},t)}$ are seventh detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{36})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{37})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{38})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$ are eighth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{44})^{(9,9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{45})^{(9,9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{46})^{(9,9,9,9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2 and 3

$$\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \begin{bmatrix} (a'_{44})^{(9)}\boxed{+(a''_{44})^{(9)}(T_{45},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{24})^{(4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{36})^{(7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8)}(T_{41},t)} \end{bmatrix} G_{13}$$

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$$\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \begin{bmatrix} (a'_{45})^{(9)}\boxed{+(a''_{45})^{(9)}(T_{45},t)} & \boxed{+(a''_{17})^{(2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{21})^{(3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{25})^{(4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{29})^{(5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{33})^{(6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{14})^{(1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{37})^{(7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{41})^{(8,8,8,8,8,8,8,8)}(T_{41},t)} \end{bmatrix} G_{14}$$

$$\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \begin{bmatrix} (a'_{46})^{(9)}\boxed{+(a''_{46})^{(9)}(T_{37},t)} & \boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{42})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)} \end{bmatrix} G_{15}$$

Where $\boxed{+(a''_{44})^{(9)}(T_{45},t)}$, $\boxed{+(a''_{45})^{(9)}(T_{45},t)}$, $\boxed{+(a''_{46})^{(9)}(T_{37},t)}$ are first augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{16})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{17})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{20})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{21})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{24})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{25})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{28})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{29})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{32})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{33})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{13})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a''_{14})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)}$ are Seventh augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)}$ $\boxed{+(a''_{37})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)}$ $\boxed{+(a''_{36})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)}$ are eighth augmentation coefficient for 1,2,3

$\boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a''_{42})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a''_{41})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)}$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{44}}{dt} = (b_{44})^{(9)}T_{45} - \begin{bmatrix} (b'_{44})^{(9)}\boxed{-(b''_{44})^{(9)}(G_{47},t)} & \boxed{-(b''_{16})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{20})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{24})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{28})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{32})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{13})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{36})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{40})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{bmatrix} T_{13}$$

$$\frac{dT_{45}}{dt} = (b_{45})^{(9)}T_{44} - \begin{bmatrix} (b'_{45})^{(9)}\boxed{-(b''_{45})^{(9)}(G_{47},t)} & \boxed{-(b''_{17})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{21})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{25})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{29})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{33})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{14})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{37})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{41})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{bmatrix} T_{14}$$

$$\frac{dT_{46}}{dt} =$$

$$(b_{46})^{(9)}T_{45} - \begin{bmatrix} (b'_{46})^{(9)}\boxed{-(b''_{46})^{(9)}(G_{47},t)} & \boxed{-(b''_{18})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{26})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{38})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{bmatrix} T_{15}$$

Where $\boxed{-(b''_{44})^{(9)}(G_{47},t)}$, $\boxed{-(b''_{45})^{(9)}(G_{47},t)}$, $\boxed{-(b''_{46})^{(9)}(G_{47},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{16})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{17})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{18})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)}$ are second detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{20})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{21})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{22})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)}$ are third detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{24})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{25})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{26})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)}$ are fourth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{28})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{29})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{30})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)}$ are fifth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{32})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b''_{33})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b''_{34})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)}$ are sixth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{15})^{(1,1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{14})^{(1,1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{13})^{(1,1,1,1,1,1,1,1,1)}(G,t)}$ are seventh detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{37})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{36})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{38})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)}$ are eighth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{42})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{41})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{40})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$ are ninth detrition coefficients for category 1, 2 and 3

Where we suppose

$(a_i)^{(1)}, (a'_i)^{(1)}, (a''_i)^{(1)}, (b_i)^{(1)}, (b'_i)^{(1)}, (b''_i)^{(1)} > 0, \quad i,j = 13,14,15$ *97*

The functions $(a''_i)^{(1)}, (b''_i)^{(1)}$ are positive continuous increasing and bounded.

**Definition of** $(p_i)^{(1)}$, $(r_i)^{(1)}$:

$$(a''_i)^{(1)}(T_{14},t) \le (p_i)^{(1)} \le (\hat{A}_{13})^{(1)}$$

$$(b''_i)^{(1)}(G,t) \le (r_i)^{(1)} \le (b'_i)^{(1)} \le (\hat{B}_{13})^{(1)}$$

$lim_{T_2\to\infty}(a''_i)^{(1)}(T_{14},t) = (p_i)^{(1)}$ 98

$\lim_{G\to\infty}(b''_i)^{(1)}(G,t) = (r_i)^{(1)}$

**Definition of** $(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}$ :

Where $\boxed{(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}}$ are positive constants and $\boxed{i = 13,14,15}$

| | |
|---|---|
| They satisfy Lipschitz condition:<br>$\|(a_i'')^{(1)}(T_{14}', t) - (a_i'')^{(1)}(T_{14}, t)\| \leq (\hat{k}_{13})^{(1)}\|T_{14} - T_{14}'\|e^{-(\hat{M}_{13})^{(1)}t}$<br>$\|(b_i'')^{(1)}(G', t) - (b_i'')^{(1)}(G, t)\| < (\hat{k}_{13})^{(1)}\|G - G'\|e^{-(\hat{M}_{13})^{(1)}t}$ | 99 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(1)}(T_{14}', t)$ and $(a_i'')^{(1)}(T_{14}, t)$ . $(T_{14}', t)$ and $(T_{14}, t)$ are points belonging to the interval $[(\hat{k}_{13})^{(1)}, (\hat{M}_{13})^{(1)}]$. It is to be noted that $(a_i'')^{(1)}(T_{14}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{13})^{(1)} = 1$ then the function $(a_i'')^{(1)}(T_{14}, t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}$ :<br>$(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)},$ are positive constants<br>$\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}} < 1$ | 100 |
| **Definition of** $(\hat{P}_{13})^{(1)}, (\hat{Q}_{13})^{(1)}$ :<br>There exists two constants $(\hat{P}_{13})^{(1)}$ and $(\hat{Q}_{13})^{(1)}$ which together With $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}, (\hat{A}_{13})^{(1)}$ and $(\hat{B}_{13})^{(1)}$ and the constants $(a_i)^{(1)}, (a_i')^{(1)}, (b_i)^{(1)}, (b_i')^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}, i = 13,14,15,$ satisfy the inequalities<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(a_i)^{(1)} + (a_i')^{(1)} + (\hat{A}_{13})^{(1)} + (\hat{P}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(b_i)^{(1)} + (b_i')^{(1)} + (\hat{B}_{13})^{(1)} + (\hat{Q}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$ | 101 |
| Where we suppose | |
| $(a_i)^{(2)}, (a_i')^{(2)}, (a_i'')^{(2)}, (b_i)^{(2)}, (b_i')^{(2)}, (b_i'')^{(2)} > 0, \quad i, j = 16,17,18$ | |
| The functions $(a_i'')^{(2)}, (b_i'')^{(2)}$ are positive continuous increasing and bounded. | |
| **Definition of** $(p_i)^{(2)}$, $(r_i)^{(2)}$: | |
| $(a_i'')^{(2)}(T_{17}, t) \leq (p_i)^{(2)} \leq (\hat{A}_{16})^{(2)}$ | 102 |
| $(b_i'')^{(2)}(G_{19}, t) \leq (r_i)^{(2)} \leq (b_i')^{(2)} \leq (\hat{B}_{16})^{(2)}$ | 103 |
| $\lim_{T_2 \to \infty}(a_i'')^{(2)}(T_{17}, t) = (p_i)^{(2)}$ | 104 |
| $\lim_{G \to \infty}(b_i'')^{(2)}((G_{19}), t) = (r_i)^{(2)}$ | 105 |
| **Definition of** $(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}$ : | 106 |

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|---|---|
| Where $\boxed{(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}, (p_i)^{(2)}, (r_i)^{(2)}}$ are positive constants and $\boxed{i = 16,17,18}$ | |
| They satisfy Lipschitz condition: | |
| $\lvert(a_i^{''})^{(2)}(T_{17}^{'}, t) - (a_i^{''})^{(2)}(T_{17}, t)\rvert \leq (\hat{k}_{16})^{(2)} \lvert T_{17} - T_{17}^{'}\rvert e^{-(\hat{M}_{16})^{(2)}t}$ | 107 |
| $\lvert(b_i^{''})^{(2)}((G_{19})^{'}, t) - (b_i^{''})^{(2)}((G_{19}), t)\rvert < (\hat{k}_{16})^{(2)} \lvert\lvert(G_{19}) - (G_{19})^{'}\rvert\rvert e^{-(\hat{M}_{16})^{(2)}t}$ | 108 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(2)}(T_{17}^{'}, t)$ and $(a_i^{''})^{(2)}(T_{17}, t)$ . $(T_{17}^{'}, t)$ and $(T_{17}, t)$ are points belonging to the interval $\left[(\hat{k}_{16})^{(2)}, (\hat{M}_{16})^{(2)}\right]$. It is to be noted that $(a_i^{''})^{(2)}(T_{17}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{16})^{(2)} = 1$ then the function $(a_i^{''})^{(2)}(T_{17}, t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$ : | |
| $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$, are positive constants<br><br>$\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}$ , $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ | 109 |
| **Definition of** $(\hat{P}_{13})^{(2)}, (\hat{Q}_{13})^{(2)}$ :<br><br>There exists two constants $(\hat{P}_{16})^{(2)}$ and $(\hat{Q}_{16})^{(2)}$ which together with $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}, (\hat{A}_{16})^{(2)} and\ (\hat{B}_{16})^{(2)}$ and the constants $(a_i)^{(2)}, (a_i^{'})^{(2)}, (b_i)^{(2)}, (b_i^{'})^{(2)}, (p_i)^{(2)}, (r_i)^{(2)}, i = 16,17,18,$<br><br>satisfy the inequalities | |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(a_i)^{(2)} + (a_i^{'})^{(2)} + (\hat{A}_{16})^{(2)} + (\hat{P}_{16})^{(2)} (\hat{k}_{16})^{(2)}] < 1$ | 110 |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(b_i)^{(2)} + (b_i^{'})^{(2)} + (\hat{B}_{16})^{(2)} + (\hat{Q}_{16})^{(2)} (\hat{k}_{16})^{(2)}] < 1$ | 111 |
| Where we suppose | |
| $(a_i)^{(3)}, (a_i^{'})^{(3)}, (a_i^{''})^{(3)}, (b_i)^{(3)}, (b_i^{'})^{(3)}, (b_i^{''})^{(3)} > 0, \quad i, j = 20,21,22$<br><br>The functions $(a_i^{''})^{(3)}, (b_i^{''})^{(3)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(3)}$, $(r_i)^{(3)}$:<br><br>$(a_i^{''})^{(3)}(T_{21}, t) \leq (p_i)^{(3)} \leq (\hat{A}_{20})^{(3)}$<br><br>$(b_i^{''})^{(3)}(G_{23}, t) \leq (r_i)^{(3)} \leq (b_i^{'})^{(3)} \leq (\hat{B}_{20})^{(3)}$ | 112 |
| $lim_{T_2 \to \infty}(a_i^{''})^{(3)}(T_{21}, t) = (p_i)^{(3)}$ | 113 |

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|---|---|
| $\lim_{G\to\infty}(b_i'')^{(3)}(G_{23},t) = (r_i)^{(3)}$<br><br>**Definition of** $(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}$ :<br><br>Where $\boxed{(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}}$ are positive constants and $\boxed{i = 20,21,22}$ | |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i'')^{(3)}(T_{21}',t) - (a_i'')^{(3)}(T_{21},t)\rvert \le (\hat{k}_{20})^{(3)}\lvert T_{21} - T_{21}'\rvert e^{-(\hat{M}_{20})^{(3)}t}$<br><br>$\lvert(b_i'')^{(3)}(G_{23}',t) - (b_i'')^{(3)}(G_{23},t)\rvert < (\hat{k}_{20})^{(3)}\lvert\lvert G_{23} - G_{23}'\rvert\rvert e^{-(\hat{M}_{20})^{(3)}t}$ | 114 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(3)}(T_{21}',t)$ and$(a_i'')^{(3)}(T_{21},t)$ . $(T_{21}',t)$ And $(T_{21},t)$ are points belonging to the interval $\left[(\hat{k}_{20})^{(3)}, (\hat{M}_{20})^{(3)}\right]$. It is to be noted that $(a_i'')^{(3)}(T_{21},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{20})^{(3)} = 1$ then the function $(a_i'')^{(3)}(T_{21},t)$, the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$ :<br><br>$(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$, are positive constants<br>$\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}$ , $\frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}} < 1$ | 115 |
| There exists two constants There exists two constants $(\hat{P}_{20})^{(3)}$ and $(\hat{Q}_{20})^{(3)}$ which together with $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}, (\hat{A}_{20})^{(3)} and\ (\hat{B}_{20})^{(3)}$ and the constants $(a_i)^{(3)}, (a_i')^{(3)}, (b_i)^{(3)}, (b_i')^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}, i = 20,21,22,$ satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(a_i)^{(3)} + (a_i')^{(3)} + (\hat{A}_{20})^{(3)} + (\hat{P}_{20})^{(3)}(\hat{k}_{20})^{(3)}] < 1$<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(b_i)^{(3)} + (b_i')^{(3)} + (\hat{B}_{20})^{(3)} + (\hat{Q}_{20})^{(3)}(\hat{k}_{20})^{(3)}] < 1$ | 116 |
| Where we suppose | |
| $(a_i)^{(4)}, (a_i')^{(4)}, (a_i'')^{(4)}, (b_i)^{(4)}, (b_i')^{(4)}, (b_i'')^{(4)} > 0, \quad i,j = 24,25,26$<br><br>The functions $(a_i'')^{(4)}, (b_i'')^{(4)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(4)}$, $(r_i)^{(4)}$:<br><br>$(a_i'')^{(4)}(T_{25},t) \le (p_i)^{(4)} \le (\hat{A}_{24})^{(4)}$<br><br>$(b_i'')^{(4)}\big((G_{27}),t\big) \le (r_i)^{(4)} \le (b_i')^{(4)} \le (\hat{B}_{24})^{(4)}$ | 117 |
| $lim_{T_2\to\infty}(a_i'')^{(4)}(T_{25},t) = (p_i)^{(4)}$ | 118 |

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|---|---|
| $\lim_{G\to\infty}(b_i'')^{(4)}\left((G_{27}),t\right) = \ (r_i)^{(4)}$<br><br>**Definition of** $(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}$ :<br><br>Where $\boxed{(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}, (p_i)^{(4)}, \ (r_i)^{(4)}}$ are positive constants and $\boxed{i = 24{,}25{,}26}$ | |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i'')^{(4)}(T_{25}',t) - (a_i'')^{(4)}(T_{25},t)\rvert \le (\hat{k}_{24})^{(4)}\lvert T_{25} - T_{25}'\rvert e^{-(\hat{M}_{24})^{(4)}t}$<br><br>$\lvert(b_i'')^{(4)}((G_{27})',t) - (b_i'')^{(4)}\left((G_{27}),t\right)\rvert < (\hat{k}_{24})^{(4)}\lvert\lvert(G_{27}) - (G_{27})'\rvert\rvert e^{-(\hat{M}_{24})^{(4)}t}$ | 119 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(4)}(T_{25}',t)$ and $(a_i'')^{(4)}(T_{25},t)$ . $(T_{25}',t)$ and $(T_{25},t)$ are points belonging to the interval $\left[(\hat{k}_{24})^{(4)}, (\hat{M}_{24})^{(4)}\right]$. It is to be noted that $(a_i'')^{(4)}(T_{25},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{24})^{(4)} = 1$ then the function $(a_i'')^{(4)}(T_{25},t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$ :<br><br>$(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$, are positive constants<br><br>$\frac{(a_i)^{(4)}}{(\hat{M}_{24})^{(4)}}$ , $\frac{(b_i)^{(4)}}{(\hat{M}_{24})^{(4)}} < 1$ | 120 |
| **Definition of** $(\hat{P}_{24})^{(4)}, (\hat{Q}_{24})^{(4)}$ :<br><br>There exists two constants $(\hat{P}_{24})^{(4)}$ and $(\hat{Q}_{24})^{(4)}$ which together with $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}, (\hat{A}_{24})^{(4)} and\ (\hat{B}_{24})^{(4)}$ and the constants $(a_i)^{(4)}, (a_i')^{(4)}, (b_i)^{(4)}, (b_i')^{(4)}, (p_i)^{(4)}, \ (r_i)^{(4)}, i = 24{,}25{,}26$, satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{24})^{(4)}}[(a_i)^{(4)} + (a_i')^{(4)} + \ (\hat{A}_{24})^{(4)} + (\hat{P}_{24})^{(4)}(\hat{k}_{24})^{(4)}] < 1$<br><br>$\frac{1}{(\hat{M}_{24})^{(4)}}[\ (b_i)^{(4)} + (b_i')^{(4)} + \ (\hat{B}_{24})^{(4)} + (\hat{Q}_{24})^{(4)}\ (\hat{k}_{24})^{(4)}] < 1$ | 121 |
| Where we suppose | |
| $(a_i)^{(5)}, (a_i')^{(5)}, (a_i'')^{(5)}, (b_i)^{(5)}, (b_i')^{(5)}, (b_i'')^{(5)} > 0, \quad i,j = 28{,}29{,}30$<br><br>The functions $(a_i'')^{(5)}, (b_i'')^{(5)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(5)}, \ (r_i)^{(5)}$:<br><br>$(a_i'')^{(5)}(T_{29},t) \le (p_i)^{(5)} \le (\hat{A}_{28})^{(5)}$<br><br>$(b_i'')^{(5)}\left((G_{31}),t\right) \le \ (r_i)^{(5)} \le (b_i')^{(5)} \le (\hat{B}_{28})^{(5)}$ | 122 |

$lim_{T_2\to\infty}(a_i'')^{(5)}(T_{29},t) = (p_i)^{(5)}$

$\lim_{G\to\infty}(b_i'')^{(5)}(G_{31},t) = (r_i)^{(5)}$

**Definition of** $(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}$ :

Where $\boxed{(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}, (p_i)^{(5)}, (r_i)^{(5)}}$ are positive constants and $\boxed{i = 28,29,30}$

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They satisfy Lipschitz condition:

$$|(a_i'')^{(5)}(T_{29}',t) - (a_i'')^{(5)}(T_{29},t)| \le (\hat{k}_{28})^{(5)}|T_{29} - T_{29}'|e^{-(\hat{M}_{28})^{(5)}t}$$

$$|(b_i'')^{(5)}((G_{31})',t) - (b_i'')^{(5)}((G_{31}),t)| < (\hat{k}_{28})^{(5)}||(G_{31}) - (G_{31})'||e^{-(\hat{M}_{28})^{(5)}t}$$

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With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(5)}(T_{29}',t)$ and $(a_i'')^{(5)}(T_{29},t)$ . $(T_{29}',t)$ and $(T_{29},t)$ are points belonging to the interval $[(\hat{k}_{28})^{(5)}, (\hat{M}_{28})^{(5)}]$. It is to be noted that $(a_i'')^{(5)}(T_{29},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{28})^{(5)} = 1$ then the function $(a_i'')^{(5)}(T_{29},t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

**Definition of** $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}$ :

$(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}$, are positive constants

$$\frac{(a_i)^{(5)}}{(\hat{M}_{28})^{(5)}}, \frac{(b_i)^{(5)}}{(\hat{M}_{28})^{(5)}} < 1$$

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**Definition of** $(\hat{P}_{28})^{(5)}, (\hat{Q}_{28})^{(5)}$ :

There exists two constants $(\hat{P}_{28})^{(5)}$ and $(\hat{Q}_{28})^{(5)}$ which together with $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}, (\hat{A}_{28})^{(5)} and\ (\hat{B}_{28})^{(5)}$ and the constants $(a_i)^{(5)}, (a_i')^{(5)}, (b_i)^{(5)}, (b_i')^{(5)}, (p_i)^{(5)}, (r_i)^{(5)}, i = 28,29,30,$ satisfy the inequalities

$$\frac{1}{(\hat{M}_{28})^{(5)}}[(a_i)^{(5)} + (a_i')^{(5)} + (\hat{A}_{28})^{(5)} + (\hat{P}_{28})^{(5)}(\hat{k}_{28})^{(5)}] < 1$$

$$\frac{1}{(\hat{M}_{28})^{(5)}}[(b_i)^{(5)} + (b_i')^{(5)} + (\hat{B}_{28})^{(5)} + (\hat{Q}_{28})^{(5)}(\hat{k}_{28})^{(5)}] < 1$$

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Where we suppose

$(a_i)^{(6)}, (a_i')^{(6)}, (a_i'')^{(6)}, (b_i)^{(6)}, (b_i')^{(6)}, (b_i'')^{(6)} > 0, \quad i,j = 32,33,34$

The functions $(a_i'')^{(6)}, (b_i'')^{(6)}$ are positive continuous increasing and bounded.

**Definition of** $(p_i)^{(6)}, (r_i)^{(6)}$:

$$(a_i'')^{(6)}(T_{33},t) \le (p_i)^{(6)} \le (\hat{A}_{32})^{(6)}$$

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| $(b_i^{''})^{(6)}((G_{35}),t) \leq \ (r_i)^{(6)} \leq (b_i^{'})^{(6)} \leq (\hat{B}_{32})^{(6)}$ | |
| $lim_{T_2\to\infty}(a_i^{''})^{(6)}(T_{33},t) = (p_i)^{(6)}$<br><br>$\lim_{G\to\infty}(b_i^{''})^{(6)}\left((G_{35}),t\right) = \ (r_i)^{(6)}$<br><br>**Definition of** $(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}$ :<br><br>Where $\boxed{(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}, (p_i)^{(6)}, \ (r_i)^{(6)}}$ are positive constants and $\boxed{i = 32,33,34}$ | 128 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(6)}(T_{33}^{'},t) - (a_i^{''})^{(6)}(T_{33},t)\rvert \leq (\hat{k}_{32})^{(6)}\lvert T_{33} - T_{33}^{'}\rvert e^{-(\hat{M}_{32})^{(6)}t}$<br><br>$\lvert(b_i^{''})^{(6)}((G_{35})^{'},t) - (b_i^{''})^{(6)}\left((G_{35}),t\right)\rvert < (\hat{k}_{32})^{(6)}\lvert\lvert(G_{35}) - (G_{35})^{'}\rvert\rvert e^{-(\hat{M}_{32})^{(6)}t}$ | |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(6)}(T_{33}^{'},t)$ and$(a_i^{''})^{(6)}(T_{33},t)$ . $(T_{33}^{'},t)$ and $(T_{33},t)$ are points belonging to the interval $\left[(\hat{k}_{32})^{(6)}, (\hat{M}_{32})^{(6)}\right]$. It is to be noted that $(a_i^{''})^{(6)}(T_{33},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{32})^{(6)} = 1$ then the function $(a_i^{''})^{(6)}(T_{33},t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}$ :<br><br>$(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}$, are positive constants<br><br>$\frac{(a_i)^{(6)}}{(\hat{M}_{32})^{(6)}}$ , $\frac{(b_i)^{(6)}}{(\hat{M}_{32})^{(6)}} < 1$ | 129 |
| **Definition of** $(\hat{P}_{32})^{(6)}, (\hat{Q}_{32})^{(6)}$ :<br><br>There exists two constants $(\hat{P}_{32})^{(6)}$ and $(\hat{Q}_{32})^{(6)}$ which together with $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}, (\hat{A}_{32})^{(6)} and\ (\hat{B}_{32})^{(6)}$ and the constants $(a_i)^{(6)}, (a_i^{'})^{(6)}, (b_i)^{(6)}, (b_i^{'})^{(6)}, (p_i)^{(6)}, \ (r_i)^{(6)}, i = 32,33,34,$ satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{32})^{(6)}}[\ (a_i)^{(6)} + (a_i^{'})^{(6)} + \ (\hat{A}_{32})^{(6)} + \ (\hat{P}_{32})^{(6)}\ (\hat{k}_{32})^{(6)}] < 1$<br><br>$\frac{1}{(\hat{M}_{32})^{(6)}}[\ (b_i)^{(6)} + (b_i^{'})^{(6)} + \ (\hat{B}_{32})^{(6)} + \ (\hat{Q}_{32})^{(6)}\ \ (\hat{k}_{32})^{(6)}] < 1$ | 130 |
| Where we suppose | |
| (M) $(a_i)^{(7)}, (a_i^{'})^{(7)}, (a_i^{''})^{(7)}, (b_i)^{(7)}, (b_i^{'})^{(7)}, (b_i^{''})^{(7)} > 0, \quad i,j = 36,37,38$<br><br>(N) The functions $(a_i^{''})^{(7)}, (b_i^{''})^{(7)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(7)}, \ (r_i)^{(7)}$: | 131 |

$(a_i^{''})^{(7)}(T_{37},t) \leq (p_i)^{(7)} \leq (\hat{A}_{36})^{(7)}$

$(b_i^{''})^{(7)}(G_{39},t) \leq (r_i)^{(7)} \leq (b_i^{'})^{(7)} \leq (\hat{B}_{36})^{(7)}$

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*(O)* $\lim_{T_2\to\infty} (a_i^{''})^{(7)}(T_{37},t) = (p_i)^{(7)}$

*(P)*

$\lim_{G\to\infty}(b_i^{''})^{(7)}\left((G_{39}),t\right) = (r_i)^{(7)}$

**Definition of** $(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}$ :

Where $\boxed{(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}, (p_i)^{(7)}, (r_i)^{(7)}}$ are positive constants and $\boxed{i = 36,37,38}$

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They satisfy Lipschitz condition:

$|(a_i^{''})^{(7)}(T_{37}^{'},t) - (a_i^{''})^{(7)}(T_{37},t)| \leq (\hat{k}_{36})^{(7)}|T_{37} - T_{37}^{'}|e^{-(\hat{M}_{36})^{(7)}t}$

$|(b_i^{''})^{(7)}((G_{39})^{'},t) - (b_i^{''})^{(7)}\left((G_{39}),t\right)| < (\hat{k}_{36})^{(7)}||(G_{39}) - (G_{39})^{'}||e^{-(\hat{M}_{36})^{(7)}t}$

With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(7)}(T_{37}^{'},t)$ and$(a_i^{''})^{(7)}(T_{37},t)$ . $(T_{37}^{'},t)$ and $(T_{37},t)$ are points belonging to the interval $\left[(\hat{k}_{36})^{(7)}, (\hat{M}_{36})^{(7)}\right]$. It is to be noted that $(a_i^{''})^{(7)}(T_{37},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{36})^{(7)} = 1$ then the function $(a_i^{''})^{(7)}(T_{37},t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

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**Definition of** $(\hat{M}_{36})^{(7)}, (\hat{k}_{36})^{(7)}$ :

(Q) $(\hat{M}_{36})^{(7)}, (\hat{k}_{36})^{(7)}$, are positive constants

$\frac{(a_i)^{(7)}}{(\hat{M}_{36})^{(7)}}, \frac{(b_i)^{(7)}}{(\hat{M}_{36})^{(7)}} < 1$

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**Definition of** $(\hat{P}_{36})^{(7)}, (\hat{Q}_{36})^{(7)}$ :

(R) There exists two constants $(\hat{P}_{36})^{(7)}$ and $(\hat{Q}_{36})^{(7)}$ which together with $(\hat{M}_{36})^{(7)}, (\hat{k}_{36})^{(7)}, (\hat{A}_{36})^{(7)} and\ (\hat{B}_{36})^{(7)}$ and the constants $(a_i)^{(7)}, (a_i^{'})^{(7)}, (b_i)^{(7)}, (b_i^{'})^{(7)}, (p_i)^{(7)}, (r_i)^{(7)}, i = 36,37,38$, satisfy the inequalities

$\frac{1}{(\hat{M}_{36})^{(7)}}[(a_i)^{(7)} + (a_i^{'})^{(7)} + (\hat{A}_{36})^{(7)} + (\hat{P}_{36})^{(7)}(\hat{k}_{36})^{(7)}] < 1$

$\frac{1}{(\hat{M}_{36})^{(7)}}[(b_i)^{(7)} + (b_i^{'})^{(7)} + (\hat{B}_{36})^{(7)} + (\hat{Q}_{36})^{(7)}(\hat{k}_{36})^{(7)}] < 1$

| | |
|---|---|
| Where we suppose | |
| $(a_i)^{(8)}, (a_i^{'})^{(8)}, (a_i^{''})^{(8)}, (b_i)^{(8)}, (b_i^{'})^{(8)}, (b_i^{''})^{(8)} > 0, \quad i,j = 40,41,42$ | 136 |
| The functions $(a_i^{''})^{(8)}, (b_i^{''})^{(8)}$ are positive continuous increasing and bounded | |
| **Definition of** $(p_i)^{(8)}$, $(r_i)^{(8)}$: | 137 |
| $(a_i^{''})^{(8)}(T_{41}, t) \leq (p_i)^{(8)} \leq (\hat{A}_{40})^{(8)}$ | 138 |
| $(b_i^{''})^{(8)}((G_{43}), t) \leq (r_i)^{(8)} \leq (b_i^{'})^{(8)} \leq (\hat{B}_{40})^{(8)}$ | 139 |
| $\lim_{T_2 \to \infty} (a_i^{''})^{(8)}(T_{41}, t) = (p_i)^{(8)}$ | 140 |
| $\lim_{G \to \infty} (b_i^{''})^{(8)}((G_{43}), t) = (r_i)^{(8)}$ | 141 |
| **Definition of** $(\hat{A}_{40})^{(8)}, (\hat{B}_{40})^{(8)}$ :<br>Where $\boxed{(\hat{A}_{40})^{(8)}, (\hat{B}_{40})^{(8)}, (p_i)^{(8)}, (r_i)^{(8)}}$ are positive constants and $\boxed{i = 40,41,42}$ | |
| They satisfy Lipschitz condition: | |
| $\lvert(a_i^{''})^{(8)}(T_{41}^{'}, t) - (a_i^{''})^{(8)}(T_{41}, t)\rvert \leq (\hat{k}_{40})^{(8)} \lvert T_{41} - T_{41}^{'}\rvert e^{-(\hat{M}_{40})^{(8)} t}$ | 142 |
| $\lvert(b_i^{''})^{(8)}((G_{43})^{'}, t) - (b_i^{''})^{(8)}((G_{43}), t)\rvert < (\hat{k}_{40})^{(8)} \lvert\lvert(G_{43}) - (G_{43})^{'}\rvert\rvert e^{-(\hat{M}_{40})^{(8)} t}$ | 143 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(8)}(T_{41}^{'}, t)$ and $(a_i^{''})^{(8)}(T_{41}, t)$ . $(T_{41}^{'}, t)$ and $(T_{41}, t)$ are points belonging to the interval $\left[(\hat{k}_{40})^{(8)}, (\hat{M}_{40})^{(8)}\right]$. It is to be noted that $(a_i^{''})^{(8)}(T_{41}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{40})^{(8)} = 1$ then the function $(a_i^{''})^{(8)}(T_{41}, t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}$ : | |
| $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}$, are positive constants | |
| $\frac{(a_i)^{(8)}}{(\hat{M}_{40})^{(8)}}$ , $\frac{(b_i)^{(8)}}{(\hat{M}_{40})^{(8)}} < 1$ | 144 |
| **Definition of** $(\hat{P}_{40})^{(8)}, (\hat{Q}_{40})^{(8)}$ :<br>There exists two constants $(\hat{P}_{40})^{(8)}$ and $(\hat{Q}_{40})^{(8)}$ which together with $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}, (\hat{A}_{40})^{(8)}$ $(\hat{B}_{40})^{(8)}$ and the constants $(a_i)^{(8)}, (a_i^{'})^{(8)}, (b_i)^{(8)}, (b_i^{'})^{(8)}, (p_i)^{(8)}, (r_i)^{(8)}, i = 40,41,42,$<br>Satisfy the inequalities | |

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| $\frac{1}{(\hat{M}_{40})^{(8)}}[(a_i)^{(8)} + (a_i')^{(8)} + (\hat{A}_{40})^{(8)} + (\hat{P}_{40})^{(8)}(\hat{k}_{40})^{(8)}] < 1$ | 145 |
| $\frac{1}{(\hat{M}_{40})^{(8)}}[(b_i)^{(8)} + (b_i')^{(8)} + (\hat{B}_{40})^{(8)} + (\hat{Q}_{40})^{(8)}(\hat{k}_{40})^{(8)}] < 1$ | 146 |
| Where we suppose | |
| $(a_i)^{(9)}, (a_i')^{(9)}, (a_i'')^{(9)}, (b_i)^{(9)}, (b_i')^{(9)}, (b_i'')^{(9)} > 0, \quad i,j = 44,45,46$<br>The functions $(a_i'')^{(9)}, (b_i'')^{(9)}$ are positive continuous increasing and bounded.<br>**Definition of** $(p_i)^{(9)}$, $(r_i)^{(9)}$:<br>$(a_i'')^{(9)}(T_{45},t) \le (p_i)^{(9)} \le (\hat{A}_{44})^{(9)}$<br>$(b_i'')^{(9)}(G_{47},t) \le (r_i)^{(9)} \le (b_i')^{(9)} \le (\hat{B}_{44})^{(9)}$ | 146 A |
| $\lim_{T_2\to\infty}(a_i'')^{(9)}(T_{45},t) = (p_i)^{(9)}$<br>$\lim_{G\to\infty}(b_i'')^{(9)}(G_{47},t) = (r_i)^{(9)}$<br>**Definition of** $(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}$ :<br>Where $\boxed{(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}, (p_i)^{(9)}, (r_i)^{(9)}}$ are positive constants and $\boxed{i = 44,45,46}$ | |
| They satisfy Lipschitz condition:<br>$\lvert(a_i'')^{(9)}(T_{45}',t) - (a_i'')^{(9)}(T_{45},t)\rvert \le (\hat{k}_{44})^{(9)}\lvert T_{45} - T_{45}'\rvert e^{-(\hat{M}_{44})^{(9)}t}$<br>$\lvert(b_i'')^{(9)}((G_{47})',t) - (b_i'')^{(9)}((G_{47}),t)\rvert < (\hat{k}_{44})^{(9)}\lvert\lvert(G_{47}) - (G_{47})'\rvert\rvert e^{-(\hat{M}_{44})^{(9)}t}$ | |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(9)}(T_{45}',t)$ and $(a_i'')^{(9)}(T_{45},t)$ . $(T_{45}',t)$ and $(T_{45},t)$ are points belonging to the interval $[(\hat{k}_{44})^{(9)}, (\hat{M}_{44})^{(9)}]$. It is to be noted that $(a_i'')^{(9)}(T_{45},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{44})^{(9)} = 1$ then the function $(a_i'')^{(9)}(T_{45},t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}$ :<br>$(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}$, are positive constants<br>$\frac{(a_i)^{(9)}}{(\hat{M}_{44})^{(9)}}, \frac{(b_i)^{(9)}}{(\hat{M}_{44})^{(9)}} < 1$ | |
| **Definition of** $(\hat{P}_{44})^{(9)}, (\hat{Q}_{44})^{(9)}$ :<br>There exists two constants $(\hat{P}_{44})^{(9)}$ and $(\hat{Q}_{44})^{(9)}$ which together with $(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}, (\hat{A}_{44})^{(9)} and (\hat{B}_{44})^{(9)}$ and the constants | |

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| $(a_i)^{(9)}, (a_i')^{(9)}, (b_i)^{(9)}, (b_i')^{(9)}, (p_i)^{(9)}, \ (r_i)^{(9)}, i = 44,45,46,$ satisfy the inequalities<br>$\frac{1}{(\hat{M}_{44})^{(9)}}[(a_i)^{(9)} + (a_i')^{(9)} + (\hat{A}_{44})^{(9)} + (\hat{P}_{44})^{(9)} (\hat{k}_{44})^{(9)}] < 1$<br>$\frac{1}{(\hat{M}_{44})^{(9)}}[(b_i)^{(9)} + (b_i')^{(9)} + (\hat{B}_{44})^{(9)} + (\hat{Q}_{44})^{(9)} (\hat{k}_{44})^{(9)}] < 1$ | |
| **Theorem 1:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>Definition of $G_i(0), T_i(0)$ :<br>$G_i(t) \le (\hat{P}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le (\hat{Q}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 147 |
| **Theorem 2 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$<br>$G_i(t) \le (\hat{P}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ , $G_i(0) = G_i^0 > 0$<br>$T_i(t) \le (\hat{Q}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ , $T_i(0) = T_i^0 > 0$ | 148 |
| **Theorem 3 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>$G_i(t) \le (\hat{P}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ , $G_i(0) = G_i^0 > 0$<br>$T_i(t) \le (\hat{Q}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ , $T_i(0) = T_i^0 > 0$ | 149 |
| **Theorem 4 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \le (\hat{P}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le (\hat{Q}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 150 |
| **Theorem 5 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \le (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le (\hat{Q}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 151 |

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| **Theorem 6 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br><br>**Definition of** $G_i(0)\,, T_i(0)$ :<br><br>$G_i(t) \le \left(\hat{P}_{32}\right)^{(6)} e^{(\hat{M}_{32})^{(6)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br><br>$T_i(t) \le \left(\hat{Q}_{32}\right)^{(6)} e^{(\hat{M}_{32})^{(6)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 152 |
| **Theorem 7:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br><br>**Definition of** $G_i(0)\,, T_i(0)$ :<br><br>$G_i(t) \le \left(\hat{P}_{36}\right)^{(7)} e^{(\hat{M}_{36})^{(7)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{36}\right)^{(7)} e^{(\hat{M}_{36})^{(7)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 |
| **Theorem 8:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br><br>**Definition of** $G_i(0)\,, T_i(0)$ :<br><br>$G_i(t) \le \left(\hat{P}_{40}\right)^{(8)} e^{(\hat{M}_{40})^{(8)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{40}\right)^{(8)} e^{(\hat{M}_{40})^{(8)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 A |
| **Theorem 9:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br><br>**Definition of** $G_i(0)\,, T_i(0)$ :<br><br>$G_i(t) \le \left(\hat{P}_{44}\right)^{(9)} e^{(\hat{M}_{44})^{(9)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br><br>$T_i(t) \le \left(\hat{Q}_{44}\right)^{(9)} e^{(\hat{M}_{44})^{(9)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 B |
| **Proof:** Consider operator $\mathcal{A}^{(1)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i$: $\mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 154 |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \le \left(\hat{P}_{13}\right)^{(1)}$ , $T_i^0 \le \left(\hat{Q}_{13}\right)^{(1)}$, | 155 |
| $0 \le G_i(t) - G_i^0 \le \left(\hat{P}_{13}\right)^{(1)} e^{(\hat{M}_{13})^{(1)}t}$ | 156 |
| $0 \le T_i(t) - T_i^0 \le \left(\hat{Q}_{13}\right)^{(1)} e^{(\hat{M}_{13})^{(1)}t}$ | 157 |
| By<br><br>$\bar{G}_{13}(t) = G_{13}^0 + \int_0^t \left[(a_{13})^{(1)} G_{14}\left(s_{(13)}\right) - \left((a_{13}')^{(1)} + a_{13}'')^{(1)}\left(T_{14}\left(s_{(13)}\right), s_{(13)}\right)\right) G_{13}\left(s_{(13)}\right)\right] ds_{(13)}$ | 158 |

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| $\bar{G}_{14}(t) = G_{14}^0 + \int_0^t \left[(a_{14})^{(1)} G_{13}(s_{(13)}) - \left((a_{14}')^{(1)} + (a_{14}'')^{(1)}(T_{14}(s_{(13)}), s_{(13)})\right) G_{14}(s_{(13)})\right] ds_{(13)}$ | |
| $\bar{G}_{15}(t) = G_{15}^0 + \int_0^t \left[(a_{15})^{(1)} G_{14}(s_{(13)}) - \left((a_{15}')^{(1)} + (a_{15}'')^{(1)}(T_{14}(s_{(13)}), s_{(13)})\right) G_{15}(s_{(13)})\right] ds_{(13)}$ | |
| $\bar{T}_{13}(t) = T_{13}^0 + \int_0^t \left[(b_{13})^{(1)} T_{14}(s_{(13)}) - \left((b_{13}')^{(1)} - (b_{13}'')^{(1)}(G(s_{(13)}), s_{(13)})\right) T_{13}(s_{(13)})\right] ds_{(13)}$ | |
| $\bar{T}_{14}(t) = T_{14}^0 + \int_0^t \left[(b_{14})^{(1)} T_{13}(s_{(13)}) - \left((b_{14}')^{(1)} - (b_{14}'')^{(1)}(G(s_{(13)}), s_{(13)})\right) T_{14}(s_{(13)})\right] ds_{(13)}$ | |
| $\bar{T}_{15}(t) = T_{15}^0 + \int_0^t \left[(b_{15})^{(1)} T_{14}(s_{(13)}) - \left((b_{15}')^{(1)} - (b_{15}'')^{(1)}(G(s_{(13)}), s_{(13)})\right) T_{15}(s_{(13)})\right] ds_{(13)}$<br>Where $s_{(13)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(2)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 159 |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{16})^{(2)}$, $T_i^0 \le (\hat{Q}_{16})^{(2)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ | |
| By<br>$\bar{G}_{16}(t) = G_{16}^0 + \int_0^t \left[(a_{16})^{(2)} G_{17}(s_{(16)}) - \left((a_{16}')^{(2)} + a_{16}'')^{(2)}(T_{17}(s_{(16)}), s_{(16)})\right) G_{16}(s_{(16)})\right] ds_{(16)}$ | 160 |
| $\bar{G}_{17}(t) = G_{17}^0 + \int_0^t \left[(a_{17})^{(2)} G_{16}(s_{(16)}) - \left((a_{17}')^{(2)} + (a_{17}'')^{(2)}(T_{17}(s_{(16)}), s_{(17)})\right) G_{17}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{G}_{18}(t) = G_{18}^0 + \int_0^t \left[(a_{18})^{(2)} G_{17}(s_{(16)}) - \left((a_{18}')^{(2)} + (a_{18}'')^{(2)}(T_{17}(s_{(16)}), s_{(16)})\right) G_{18}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{T}_{16}(t) = T_{16}^0 + \int_0^t \left[(b_{16})^{(2)} T_{17}(s_{(16)}) - \left((b_{16}')^{(2)} - (b_{16}'')^{(2)}(G_{19}(s_{(16)}), s_{(16)})\right) T_{16}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{T}_{17}(t) = T_{17}^0 + \int_0^t \left[(b_{17})^{(2)} T_{16}(s_{(16)}) - \left((b_{17}')^{(2)} - (b_{17}'')^{(2)}(G_{19}(s_{(16)}), s_{(16)})\right) T_{17}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{T}_{18}(t) = T_{18}^0 + \int_0^t \left[(b_{18})^{(2)} T_{17}(s_{(16)}) - \left((b_{18}')^{(2)} - (b_{18}'')^{(2)}(G_{19}(s_{(16)}), s_{(16)})\right) T_{18}(s_{(16)})\right] ds_{(16)}$<br>Where $s_{(16)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(3)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |

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| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{20})^{(3)}$ , $T_i^0 \leq (\hat{Q}_{20})^{(3)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{20})^{(3)} e^{(\bar{M}_{20})^{(3)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{20})^{(3)} e^{(\bar{M}_{20})^{(3)} t}$ | |
| By<br>$\bar{G}_{20}(t) = G_{20}^0 + \int_0^t \left[ (a_{20})^{(3)} G_{21}(s_{(20)}) - \left( (a'_{20})^{(3)} + a''_{20})^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{20}(s_{(20)}) \right] ds_{(20)}$ | 161 |
| $\bar{G}_{21}(t) = G_{21}^0 + \int_0^t \left[ (a_{21})^{(3)} G_{20}(s_{(20)}) - \left( (a'_{21})^{(3)} + (a''_{21})^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{21}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{G}_{22}(t) = G_{22}^0 + \int_0^t \left[ (a_{22})^{(3)} G_{21}(s_{(20)}) - \left( (a'_{22})^{(3)} + (a''_{22})^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{22}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{20}(t) = T_{20}^0 + \int_0^t \left[ (b_{20})^{(3)} T_{21}(s_{(20)}) - \left( (b'_{20})^{(3)} - (b''_{20})^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{20}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{21}(t) = T_{21}^0 + \int_0^t \left[ (b_{21})^{(3)} T_{20}(s_{(20)}) - \left( (b'_{21})^{(3)} - (b''_{21})^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{21}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{22}(t) = T_{22}^0 + \int_0^t \left[ (b_{22})^{(3)} T_{21}(s_{(20)}) - \left( (b'_{22})^{(3)} - (b''_{22})^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{22}(s_{(20)}) \right] ds_{(20)}$<br>Where $s_{(20)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:** Consider operator $\mathcal{A}^{(4)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i$: $\mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{24})^{(4)}$ , $T_i^0 \leq (\hat{Q}_{24})^{(4)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{24})^{(4)} e^{(\bar{M}_{24})^{(4)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{24})^{(4)} e^{(\bar{M}_{24})^{(4)} t}$ | |
| By<br>$\bar{G}_{24}(t) = G_{24}^0 + \int_0^t \left[ (a_{24})^{(4)} G_{25}(s_{(24)}) - \left( (a'_{24})^{(4)} + a''_{24})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{24}(s_{(24)}) \right] ds_{(24)}$ | 162 |
| $\bar{G}_{25}(t) = G_{25}^0 + \int_0^t \left[ (a_{25})^{(4)} G_{24}(s_{(24)}) - \left( (a'_{25})^{(4)} + (a''_{25})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{25}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{G}_{26}(t) = G_{26}^0 + \int_0^t \left[ (a_{26})^{(4)} G_{25}(s_{(24)}) - \left( (a'_{26})^{(4)} + (a''_{26})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{26}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{24}(t) = T_{24}^0 + \int_0^t \left[ (b_{24})^{(4)} T_{25}(s_{(24)}) - \left( (b'_{24})^{(4)} - (b''_{24})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{24}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{25}(t) = T_{25}^0 + \int_0^t \left[ (b_{25})^{(4)} T_{24}(s_{(24)}) - \left( (b'_{25})^{(4)} - (b''_{25})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{25}(s_{(24)}) \right] ds_{(24)}$ | |

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| $\overline{T}_{26}(t) = T_{26}^0 + \int_0^t \left[(b_{26})^{(4)} T_{25}(s_{(24)}) - \left((b_{26}')^{(4)} - (b_{26}'')^{(4)}(G_{27}(s_{(24)}), s_{(24)})\right) T_{26}(s_{(24)})\right] ds_{(24)}$<br>Where $s_{(24)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:** Consider operator $\mathcal{A}^{(5)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{28})^{(5)}$ , $T_i^0 \leq (\hat{Q}_{28})^{(5)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| By<br>$\bar{G}_{28}(t) = G_{28}^0 + \int_0^t \left[(a_{28})^{(5)} G_{29}(s_{(28)}) - \left((a_{28}')^{(5)} + a_{28}'')^{(5)}(T_{29}(s_{(28)}), s_{(28)})\right) G_{28}(s_{(28)})\right] ds_{(28)}$ | 163 |
| $\bar{G}_{29}(t) = G_{29}^0 + \int_0^t \left[(a_{29})^{(5)} G_{28}(s_{(28)}) - \left((a_{29}')^{(5)} + (a_{29}'')^{(5)}(T_{29}(s_{(28)}), s_{(28)})\right) G_{29}(s_{(28)})\right] ds_{(28)}$ | |
| $\bar{G}_{30}(t) = G_{30}^0 + \int_0^t \left[(a_{30})^{(5)} G_{29}(s_{(28)}) - \left((a_{30}')^{(5)} + (a_{30}'')^{(5)}(T_{29}(s_{(28)}), s_{(28)})\right) G_{30}(s_{(28)})\right] ds_{(28)}$ | |
| $\overline{T}_{28}(t) = T_{28}^0 + \int_0^t \left[(b_{28})^{(5)} T_{29}(s_{(28)}) - \left((b_{28}')^{(5)} - (b_{28}'')^{(5)}(G_{31}(s_{(28)}), s_{(28)})\right) T_{28}(s_{(28)})\right] ds_{(28)}$ | |
| $\overline{T}_{29}(t) = T_{29}^0 + \int_0^t \left[(b_{29})^{(5)} T_{28}(s_{(28)}) - \left((b_{29}')^{(5)} - (b_{29}'')^{(5)}(G_{31}(s_{(28)}), s_{(28)})\right) T_{29}(s_{(28)})\right] ds_{(28)}$ | |
| $\overline{T}_{30}(t) = T_{30}^0 + \int_0^t \left[(b_{30})^{(5)} T_{29}(s_{(28)}) - \left((b_{30}')^{(5)} - (b_{30}'')^{(5)}(G_{31}(s_{(28)}), s_{(28)})\right) T_{30}(s_{(28)})\right] ds_{(28)}$<br>Where $s_{(28)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(6)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{32})^{(6)}$ , $T_i^0 \leq (\hat{Q}_{32})^{(6)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ | |
| By<br>$\bar{G}_{32}(t) = G_{32}^0 + \int_0^t \left[(a_{32})^{(6)} G_{33}(s_{(32)}) - \left((a_{32}')^{(6)} + a_{32}'')^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{32}(s_{(32)})\right] ds_{(32)}$ | 164 |
| $\bar{G}_{33}(t) = G_{33}^0 + \int_0^t \left[(a_{33})^{(6)} G_{32}(s_{(32)}) - \left((a_{33}')^{(6)} + (a_{33}'')^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{33}(s_{(32)})\right] ds_{(32)}$ | |

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| $\bar{G}_{34}(t) = G_{34}^0 + \int_0^t \left[(a_{34})^{(6)} G_{33}(s_{(32)}) - \left((a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{34}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{T}_{32}(t) = T_{32}^0 + \int_0^t \left[(b_{32})^{(6)} T_{33}(s_{(32)}) - \left((b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{32}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{T}_{33}(t) = T_{33}^0 + \int_0^t \left[(b_{33})^{(6)} T_{32}(s_{(32)}) - \left((b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{33}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{T}_{34}(t) = T_{34}^0 + \int_0^t \left[(b_{34})^{(6)} T_{33}(s_{(32)}) - \left((b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{34}(s_{(32)})\right] ds_{(32)}$<br>Where $s_{(32)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(7)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{36})^{(7)}$, $T_i^0 \le (\hat{Q}_{36})^{(7)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ | |
| By<br>$\bar{G}_{36}(t) = G_{36}^0 + \int_0^t \left[(a_{36})^{(7)} G_{37}(s_{(36)}) - \left((a_{36}^{'})^{(7)} + a_{36}^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{36}(s_{(36)})\right] ds_{(36)}$ | 165 |
| $\bar{G}_{37}(t) = G_{37}^0 + \int_0^t \left[(a_{37})^{(7)} G_{36}(s_{(36)}) - \left((a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{37}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{G}_{38}(t) = G_{38}^0 + \int_0^t \left[(a_{38})^{(7)} G_{37}(s_{(36)}) - \left((a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{38}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{36}(t) = T_{36}^0 + \int_0^t \left[(b_{36})^{(7)} T_{37}(s_{(36)}) - \left((b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{36}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{37}(t) = T_{37}^0 + \int_0^t \left[(b_{37})^{(7)} T_{36}(s_{(36)}) - \left((b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{37}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{38}(t) = T_{38}^0 + \int_0^t \left[(b_{38})^{(7)} T_{37}(s_{(36)}) - \left((b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{38}(s_{(36)})\right] ds_{(36)}$<br>Where $s_{(36)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(8)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ | |

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| which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{40})^{(8)}$ , $T_i^0 \leq (\hat{Q}_{40})^{(8)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{40})^{(8)} e^{(\bar{M}_{40})^{(8)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{40})^{(8)} e^{(\bar{M}_{40})^{(8)} t}$ | |
| By<br>$\bar{G}_{40}(t) = G_{40}^0 + \int_0^t \left[ (a_{40})^{(8)} G_{41}(s_{(40)}) - \left( (a'_{40})^{(8)} + a''_{40})^{(8)} (T_{41}(s_{(40)}), s_{(40)}) \right) G_{40}(s_{(40)}) \right] ds_{(40)}$ | 166 |
| $\bar{G}_{41}(t) = G_{41}^0 + \int_0^t \left[ (a_{41})^{(8)} G_{40}(s_{(40)}) - \left( (a'_{41})^{(8)} + (a''_{41})^{(8)} (T_{41}(s_{(40)}), s_{(40)}) \right) G_{41}(s_{(40)}) \right] ds_{(40)}$ | |
| $\bar{G}_{42}(t) = G_{42}^0 + \int_0^t \left[ (a_{42})^{(8)} G_{41}(s_{(40)}) - \left( (a'_{42})^{(8)} + (a''_{42})^{(8)} (T_{41}(s_{(40)}), s_{(40)}) \right) G_{42}(s_{(40)}) \right] ds_{(40)}$ | |
| $\bar{T}_{40}(t) = T_{40}^0 + \int_0^t \left[ (b_{40})^{(8)} T_{41}(s_{(40)}) - \left( (b'_{40})^{(8)} - (b''_{40})^{(8)} (G_{43}(s_{(40)}), s_{(40)}) \right) T_{40}(s_{(40)}) \right] ds_{(40)}$ | |
| $\bar{T}_{41}(t) = T_{41}^0 + \int_0^t \left[ (b_{41})^{(8)} T_{40}(s_{(40)}) - \left( (b'_{41})^{(8)} - (b''_{41})^{(8)} (G_{43}(s_{(40)}), s_{(40)}) \right) T_{41}(s_{(40)}) \right] ds_{(40)}$ | |
| $\bar{T}_{42}(t) = T_{42}^0 + \int_0^t \left[ (b_{42})^{(8)} T_{41}(s_{(40)}) - \left( (b'_{42})^{(8)} - (b''_{42})^{(8)} (G_{43}(s_{(40)}), s_{(40)}) \right) T_{42}(s_{(40)}) \right] ds_{(40)}$<br>Where $s_{(40)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(9)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 166 A |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{44})^{(9)}$ , $T_i^0 \leq (\hat{Q}_{44})^{(9)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{44})^{(9)} e^{(\bar{M}_{44})^{(9)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{44})^{(9)} e^{(\bar{M}_{44})^{(9)} t}$ | |
| By<br>$\bar{G}_{44}(t) = G_{44}^0 + \int_0^t \left[ (a_{44})^{(9)} G_{45}(s_{(44)}) - \left( (a'_{44})^{(9)} + a''_{44})^{(9)} (T_{45}(s_{(44)}), s_{(44)}) \right) G_{44}(s_{(44)}) \right] ds_{(44)}$ | |
| $\bar{G}_{45}(t) = G_{45}^0 + \int_0^t \left[ (a_{45})^{(9)} G_{44}(s_{(44)}) - \left( (a'_{45})^{(9)} + (a''_{45})^{(9)} (T_{45}(s_{(44)}), s_{(44)}) \right) G_{45}(s_{(44)}) \right] ds_{(44)}$ | |
| $\bar{G}_{46}(t) = G_{46}^0 + \int_0^t \left[ (a_{46})^{(9)} G_{45}(s_{(44)}) - \left( (a'_{46})^{(9)} + (a''_{46})^{(9)} (T_{45}(s_{(44)}), s_{(44)}) \right) G_{46}(s_{(44)}) \right] ds_{(44)}$ | |

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| $\bar{T}_{44}(t) = T_{44}^0 + \int_0^t \left[(b_{44})^{(9)} T_{45}(s_{(44)}) - \left((b'_{44})^{(9)} - (b''_{44})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{44}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{45}(t) = T_{45}^0 + \int_0^t \left[(b_{45})^{(9)} T_{44}(s_{(44)}) - \left((b'_{45})^{(9)} - (b''_{45})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{45}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{46}(t) = T_{46}^0 + \int_0^t \left[(b_{46})^{(9)} T_{45}(s_{(44)}) - \left((b'_{46})^{(9)} - (b''_{46})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{46}(s_{(44)})\right] ds_{(44)}$<br>Where $s_{(44)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| The operator $\mathcal{A}^{(1)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{13}(t) \le G_{13}^0 + \int_0^t \left[(a_{13})^{(1)} \left(G_{14}^0 + (\hat{P}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} s_{(13)}}\right)\right] ds_{(13)} =$<br>$\left(1 + (a_{13})^{(1)} t\right) G_{14}^0 + \frac{(a_{13})^{(1)} (\hat{P}_{13})^{(1)}}{(\hat{M}_{13})^{(1)}} \left(e^{(\hat{M}_{13})^{(1)} t} - 1\right)$ | 167 |
| From which it follows that<br>$(G_{13}(t) - G_{13}^0) e^{-(\hat{M}_{13})^{(1)} t} \le \frac{(a_{13})^{(1)}}{(\hat{M}_{13})^{(1)}} \left[\left((\hat{P}_{13})^{(1)} + G_{14}^0\right) e^{\left(-\frac{(\hat{P}_{13})^{(1)} + G_{14}^0}{G_{14}^0}\right)} + (\hat{P}_{13})^{(1)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 1 | 168 |
| Analogous inequalities hold also for $G_{14}, G_{15}, T_{13}, T_{14}, T_{15}$ | |
| The operator $\mathcal{A}^{(2)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{16}(t) \le G_{16}^0 + \int_0^t \left[(a_{16})^{(2)} \left(G_{17}^0 + (\hat{P}_{16})^{(6)} e^{(\hat{M}_{16})^{(2)} s_{(16)}}\right)\right] ds_{(16)} =$<br>$\left(1 + (a_{16})^{(2)} t\right) G_{17}^0 + \frac{(a_{16})^{(2)} (\hat{P}_{16})^{(2)}}{(\hat{M}_{16})^{(2)}} \left(e^{(\hat{M}_{16})^{(2)} t} - 1\right)$ | 169 |
| From which it follows that<br>$(G_{16}(t) - G_{16}^0) e^{-(\hat{M}_{16})^{(2)} t} \le \frac{(a_{16})^{(2)}}{(\hat{M}_{16})^{(2)}} \left[\left((\hat{P}_{16})^{(2)} + G_{17}^0\right) e^{\left(-\frac{(\hat{P}_{16})^{(2)} + G_{17}^0}{G_{17}^0}\right)} + (\hat{P}_{16})^{(2)}\right]$ | 170 |
| Analogous inequalities hold also for $G_{17}, G_{18}, T_{16}, T_{17}, T_{18}$ | |
| The operator $\mathcal{A}^{(3)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{20}(t) \le G_{20}^0 + \int_0^t \left[(a_{20})^{(3)} \left(G_{21}^0 + (\hat{P}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} s_{(20)}}\right)\right] ds_{(20)} =$<br>$\left(1 + (a_{20})^{(3)} t\right) G_{21}^0 + \frac{(a_{20})^{(3)} (\hat{P}_{20})^{(3)}}{(\hat{M}_{20})^{(3)}} \left(e^{(\hat{M}_{20})^{(3)} t} - 1\right)$ | 171 |
| From which it follows that | 172 |

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| $(G_{20}(t)-G_{20}^0)e^{-(\hat{M}_{20})^{(3)}t} \le \frac{(a_{20})^{(3)}}{(\hat{M}_{20})^{(3)}}\left[((\hat{P}_{20})^{(3)}+G_{21}^0)e^{\left(-\frac{(\hat{P}_{20})^{(3)}+G_{21}^0}{G_{21}^0}\right)}+(\hat{P}_{20})^{(3)}\right]$ | |
| Analogous inequalities hold also for $G_{21}, G_{22}, T_{20}, T_{21}, T_{22}$ | |
| The operator $\mathcal{A}^{(4)}$ maps the space of functions satisfying into itself .Indeed it is obvious that $G_{24}(t) \le G_{24}^0 + \int_0^t \left[(a_{24})^{(4)}\left(G_{25}^0+(\hat{P}_{24})^{(4)}e^{(\hat{M}_{24})^{(4)}s_{(24)}}\right)\right] ds_{(24)} =$ $\left(1+(a_{24})^{(4)}t\right)G_{25}^0 + \frac{(a_{24})^{(4)}(\hat{P}_{24})^{(4)}}{(\hat{M}_{24})^{(4)}}\left(e^{(\hat{M}_{24})^{(4)}t}-1\right)$ | 173 |
| From which it follows that $(G_{24}(t)-G_{24}^0)e^{-(\hat{M}_{24})^{(4)}t} \le \frac{(a_{24})^{(4)}}{(\hat{M}_{24})^{(4)}}\left[((\hat{P}_{24})^{(4)}+G_{25}^0)e^{\left(-\frac{(\hat{P}_{24})^{(4)}+G_{25}^0}{G_{25}^0}\right)}+(\hat{P}_{24})^{(4)}\right]$ $(G_i^0)$ is as defined in the statement of theorem 4 | 174 |
| The operator $\mathcal{A}^{(5)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{28}(t) \le G_{28}^0 + \int_0^t \left[(a_{28})^{(5)}\left(G_{29}^0+(\hat{P}_{28})^{(5)}e^{(\hat{M}_{28})^{(5)}s_{(28)}}\right)\right] ds_{(28)} =$ $\left(1+(a_{28})^{(5)}t\right)G_{29}^0 + \frac{(a_{28})^{(5)}(\hat{P}_{28})^{(5)}}{(\hat{M}_{28})^{(5)}}\left(e^{(\hat{M}_{28})^{(5)}t}-1\right)$ | |
| From which it follows that $(G_{28}(t)-G_{28}^0)e^{-(\hat{M}_{28})^{(5)}t} \le \frac{(a_{28})^{(5)}}{(\hat{M}_{28})^{(5)}}\left[((\hat{P}_{28})^{(5)}+G_{29}^0)e^{\left(-\frac{(\hat{P}_{28})^{(5)}+G_{29}^0}{G_{29}^0}\right)}+(\hat{P}_{28})^{(5)}\right]$ $(G_i^0)$ is as defined in the statement of theorem 5 | 175 |
| The operator $\mathcal{A}^{(6)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{32}(t) \le G_{32}^0 + \int_0^t \left[(a_{32})^{(6)}\left(G_{33}^0+(\hat{P}_{32})^{(6)}e^{(\hat{M}_{32})^{(6)}s_{(32)}}\right)\right] ds_{(32)} =$ $\left(1+(a_{32})^{(6)}t\right)G_{33}^0 + \frac{(a_{32})^{(6)}(\hat{P}_{32})^{(6)}}{(\hat{M}_{32})^{(6)}}\left(e^{(\hat{M}_{32})^{(6)}t}-1\right)$ | 176 |
| From which it follows that $(G_{32}(t)-G_{32}^0)e^{-(\hat{M}_{32})^{(6)}t} \le \frac{(a_{32})^{(6)}}{(\hat{M}_{32})^{(6)}}\left[((\hat{P}_{32})^{(6)}+G_{33}^0)e^{\left(-\frac{(\hat{P}_{32})^{(6)}+G_{33}^0}{G_{33}^0}\right)}+(\hat{P}_{32})^{(6)}\right]$ $(G_i^0)$ is as defined in the statement of theorem 6 | 177 |

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| Analogous inequalities hold also for $G_{25}, G_{26}, T_{24}, T_{25}, T_{26}$ | |
| (c) The operator $\mathcal{A}^{(7)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{36}(t) \le G_{36}^0 + \int_0^t \left[(a_{36})^{(7)} \left(G_{37}^0 + (\hat{P}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} s_{(36)}}\right)\right] ds_{(36)} =$<br>$\left(1 + (a_{36})^{(7)} t\right) G_{37}^0 + \frac{(a_{36})^{(7)} (\hat{P}_{36})^{(7)}}{(\hat{M}_{36})^{(7)}} \left(e^{(\hat{M}_{36})^{(7)} t} - 1\right)$ | 178 |
| From which it follows that<br>$(G_{36}(t) - G_{36}^0) e^{-(\hat{M}_{36})^{(7)} t} \le \frac{(a_{36})^{(7)}}{(\hat{M}_{36})^{(7)}} \left[\left((\hat{P}_{36})^{(7)} + G_{37}^0\right) e^{\left(-\frac{(\hat{P}_{36})^{(7)} + G_{37}^0}{G_{37}^0}\right)} + (\hat{P}_{36})^{(7)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 7 | |
| The operator $\mathcal{A}^{(8)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that | |
| $G_{40}(t) \le G_{40}^0 + \int_0^t \left[(a_{40})^{(8)} \left(G_{41}^0 + (\hat{P}_{40})^{(8)} e^{(\hat{M}_{40})^{(8)} s_{(40)}}\right)\right] ds_{(40)} =$<br>$\left(1 + (a_{40})^{(8)} t\right) G_{41}^0 + \frac{(a_{40})^{(8)} (\hat{P}_{40})^{(8)}}{(\hat{M}_{40})^{(8)}} \left(e^{(\hat{M}_{40})^{(8)} t} - 1\right)$ | 180 |
| From which it follows that<br>$(G_{40}(t) - G_{40}^0) e^{-(\hat{M}_{40})^{(8)} t} \le \frac{(a_{40})^{(8)}}{(\hat{M}_{40})^{(8)}} \left[\left((\hat{P}_{40})^{(8)} + G_{41}^0\right) e^{\left(-\frac{(\hat{P}_{40})^{(8)} + G_{41}^0}{G_{41}^0}\right)} + (\hat{P}_{40})^{(8)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 8<br>Analogous inequalities hold also for $G_{41}, G_{42}, T_{40}, T_{41}, T_{42}$ | 181 |
| The operator $\mathcal{A}^{(9)}$ maps the space of functions satisfying 34,35,36 into itself .Indeed it is obvious that<br>$G_{44}(t) \le G_{44}^0 + \int_0^t \left[(a_{44})^{(9)} \left(G_{45}^0 + (\hat{P}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} s_{(44)}}\right)\right] ds_{(44)} =$<br>$\left(1 + (a_{44})^{(9)} t\right) G_{45}^0 + \frac{(a_{44})^{(9)} (\hat{P}_{44})^{(9)}}{(\hat{M}_{44})^{(9)}} \left(e^{(\hat{M}_{44})^{(9)} t} - 1\right)$ | |
| From which it follows that<br>$(G_{44}(t) - G_{44}^0) e^{-(\hat{M}_{44})^{(9)} t} \le \frac{(a_{44})^{(9)}}{(\hat{M}_{44})^{(9)}} \left[\left((\hat{P}_{44})^{(9)} + G_{45}^0\right) e^{\left(-\frac{(\hat{P}_{44})^{(9)} + G_{45}^0}{G_{45}^0}\right)} + (\hat{P}_{44})^{(9)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 9<br>Analogous inequalities hold also for $G_{45}, G_{46}, T_{44}, T_{45}, T_{46}$ | |
| It is now sufficient to take $\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}} < 1$ and to choose | 182 |

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| $(\hat{P}_{13})^{(1)}$ and $(\hat{Q}_{13})^{(1)}$ large to have | |
| $\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}\left[(\hat{P}_{13})^{(1)}+\left((\hat{P}_{13})^{(1)}+G_j^0\right)e^{-\left(\frac{(\hat{P}_{13})^{(1)}+G_j^0}{G_j^0}\right)}\right]\leq(\hat{P}_{13})^{(1)}$ | 183 |
| $\frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}}\left[\left((\hat{Q}_{13})^{(1)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{13})^{(1)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{13})^{(1)}\right]\leq(\hat{Q}_{13})^{(1)}$ | 184 |
| In order that the operator $\mathcal{A}^{(1)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(1)}$ is a contraction with respect to the metric<br>$d\left(\left(G^{(1)},T^{(1)}\right),\left(G^{(2)},T^{(2)}\right)\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\hat{M}_{13})^{(1)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\hat{M}_{13})^{(1)}t}\}$ | 185 |
| Indeed if we denote<br>**Definition of** $\tilde{G},\tilde{T}:\left(\tilde{G},\tilde{T}\right)=\mathcal{A}^{(1)}(G,T)$<br>It results<br>$\left\|\tilde{G}_{13}^{(1)}-\tilde{G}_i^{(2)}\right\|\leq\int_0^t(a_{13})^{(1)}\left\|G_{14}^{(1)}-G_{14}^{(2)}\right\|e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{(\hat{M}_{13})^{(1)}s_{(13)}}ds_{(13)}+$<br>$\int_0^t\{(a_{13}')^{(1)}\left\|G_{13}^{(1)}-G_{13}^{(2)}\right\|e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{-(\hat{M}_{13})^{(1)}s_{(13)}}+$<br>$(a_{13}'')^{(1)}\left(T_{14}^{(1)},s_{(13)}\right)\left\|G_{13}^{(1)}-G_{13}^{(2)}\right\|e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{(\hat{M}_{13})^{(1)}s_{(13)}}+$<br>$G_{13}^{(2)}\left\\|(a_{13}'')^{(1)}\left(T_{14}^{(1)},s_{(13)}\right)-(a_{13}'')^{(1)}\left(T_{14}^{(2)},s_{(13)}\right)\right\\|\ e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{(\hat{M}_{13})^{(1)}s_{(13)}}\}ds_{(13)}$<br>Where $s_{(13)}$ represents integrand that is integrated over the interval $[0,t]$<br>From the hypotheses it follows | |
| $\left\|G^{(1)}-G^{(2)}\right\|e^{-(\hat{M}_{13})^{(1)}t}\leq$<br>$\frac{1}{(\hat{M}_{13})^{(1)}}\left((a_{13})^{(1)}+(a_{13}')^{(1)}+(\hat{A}_{13})^{(1)}+(\hat{P}_{13})^{(1)}(\hat{k}_{13})^{(1)}\right)d\left(\left(G^{(1)},T^{(1)};\ G^{(2)},T^{(2)}\right)\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 186 |
| **Remark 1:** The fact that we supposed $(a_{13}'')^{(1)}$ and $(b_{13}'')^{(1)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by | |

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| $(\widehat{P}_{13})^{(1)}e^{(\widehat{M}_{13})^{(1)}t}$ *and* $(\widehat{Q}_{13})^{(1)}e^{(\widehat{M}_{13})^{(1)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(1)}$ and $(b_i^{''})^{(1)}, i = 13,14,15$ depend only on $T_{14}$ and respectively on $G(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 2:** There does not exist any $t$ where $G_i(t) = 0\ and\ T_i(t) = 0$<br><br>From 19 to 24 it results<br><br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(1)} - (a_i^{''})^{(1)}(T_{14}(s_{(13)}), s_{(13)})\} ds_{(13)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(1)}t\right)} > 0$ for t > 0 | |
| **Definition of** $\left((\widehat{M}_{13})^{(1)}\right)_1, \left((\widehat{M}_{13})^{(1)}\right)_2\ and\ \left((\widehat{M}_{13})^{(1)}\right)_3$ :<br><br>**Remark 3**: if $G_{13}$ is bounded, the same property have also $G_{14}\ and\ G_{15}$ . indeed if<br><br>$G_{13} < (\widehat{M}_{13})^{(1)}$ it follows $\frac{dG_{14}}{dt} \leq \left((\widehat{M}_{13})^{(1)}\right)_1 - (a_{14}^{'})^{(1)}G_{14}$ and by integrating<br><br>$G_{14} \leq \left((\widehat{M}_{13})^{(1)}\right)_2 = G_{14}^0 + 2(a_{14})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_1/(a_{14}^{'})^{(1)}$<br><br>In the same way , one can obtain<br><br>$G_{15} \leq \left((\widehat{M}_{13})^{(1)}\right)_3 = G_{15}^0 + 2(a_{15})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_2/(a_{15}^{'})^{(1)}$<br><br>If $G_{14}\ or\ G_{15}$ is bounded, the same property follows for $G_{13}$ , $G_{15}$ and $G_{13}$ , $G_{14}$ respectively. | 187 |
| **Remark 4:** If $G_{13}$ *is* bounded, from below, the same property holds for $G_{14}\ and\ G_{15}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{14}$ is bounded from below. | 188 |
| **Remark 5:** If $T_{13}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(1)}(G(t), t)) = (b_{14}^{'})^{(1)}$ then $T_{14} \to \infty$.<br><br>**Definition of** $(m)^{(1)}$ and $\varepsilon_1$ :<br><br>Indeed let $t_1$ be so that for $t > t_1$<br><br>$(b_{14})^{(1)} - (b_i^{''})^{(1)}(G(t), t) < \varepsilon_1, T_{13}(t) > (m)^{(1)}$ | 189 |
| Then $\frac{dT_{14}}{dt} \geq (a_{14})^{(1)}(m)^{(1)} - \varepsilon_1 T_{14}$ which leads to<br><br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{\varepsilon_1}\right)(1 - e^{-\varepsilon_1 t}) + T_{14}^0 e^{-\varepsilon_1 t}$ If we take t such that $e^{-\varepsilon_1 t} = \frac{1}{2}$ it results<br><br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{2}\right),\quad t = log\frac{2}{\varepsilon_1}$ By taking now $\varepsilon_1$ sufficiently small one sees that $T_{14}$ is unbounded.<br>The same property holds for $T_{15}$ if $\lim_{t\to\infty}(b_{15}^{''})^{(1)}(G(t), t) = (b_{15}^{'})^{(1)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |

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| It is now sufficient to take $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}$ , $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ and to choose<br>$(\hat{P}_{16})^{(2)}$ *and* $(\hat{Q}_{16})^{(2)}$ large to have | 190 |
| $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[(\hat{P}_{16})^{(2)} + \left((\hat{P}_{16})^{(2)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{16})^{(2)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{16})^{(2)}$ | 191 |
| $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[\left((\hat{Q}_{16})^{(2)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{16})^{(2)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{16})^{(2)}\right] \leq (\hat{Q}_{16})^{(2)}$ | 192 |
| In order that the operator $\mathcal{A}^{(2)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 193 |
| The operator $\mathcal{A}^{(2)}$ is a contraction with respect to the metric<br>$d\left(((G_{19})^{(1)}, (T_{19})^{(1)}), ((G_{19})^{(2)}, (T_{19})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\hat{M}_{16})^{(2)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\hat{M}_{16})^{(2)}t}\}$ | 194 |
| Indeed if we denote<br>**Definition of** $\widetilde{G_{19}}, \widetilde{T_{19}}$ : $\left(\widetilde{G_{19}}, \widetilde{T_{19}}\right) = \mathcal{A}^{(2)}(G_{19}, T_{19})$ | 195 |
| It results<br>$\left\|\tilde{G}_{16}^{(1)} - \tilde{G}_i^{(2)}\right\| \leq \int_0^t (a_{16})^{(2)}\left\|G_{17}^{(1)} - G_{17}^{(2)}\right\|e^{-(\hat{M}_{16})^{(2)}s_{(16)}}e^{(\hat{M}_{16})^{(2)}s_{(16)}}\,ds_{(16)} +$<br>$\int_0^t\{(a_{16}')^{(2)}\left\|G_{16}^{(1)} - G_{16}^{(2)}\right\|e^{-(\hat{M}_{16})^{(2)}s_{(16)}}e^{-(\hat{M}_{16})^{(2)}s_{(16)}} +$<br>$(a_{16}'')^{(2)}\left(T_{17}^{(1)}, s_{(16)}\right)\left\|G_{16}^{(1)} - G_{16}^{(2)}\right\|e^{-(\hat{M}_{16})^{(2)}s_{(16)}}e^{(\hat{M}_{16})^{(2)}s_{(16)}} +$<br>$G_{16}^{(2)}\left\|(a_{16}'')^{(2)}\left(T_{17}^{(1)}, s_{(16)}\right) - (a_{16}'')^{(2)}\left(T_{17}^{(2)}, s_{(16)}\right)\right\| \; e^{-(\hat{M}_{16})^{(2)}s_{(16)}}e^{(\hat{M}_{16})^{(2)}s_{(16)}}\}ds_{(16)}$ | 196 |
| Where $s_{(16)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | 197 |
| $\left\|(G_{19})^{(1)} - (G_{19})^{(2)}\right\|e^{-(\hat{M}_{16})^{(2)}t} \leq$<br>$\frac{1}{(\hat{M}_{16})^{(2)}}\left((a_{16})^{(2)} + (a_{16}')^{(2)} + (\hat{A}_{16})^{(2)} + (\hat{P}_{16})^{(2)}(\hat{k}_{16})^{(2)}\right)d\left(((G_{19})^{(1)}, (T_{19})^{(1)}; (G_{19})^{(2)}, (T_{19})^{(2)})\right)$ | |
| And analogous inequalities for $G_i$ and $T_i$. Taking into account the hypothesis the result follows | 198 |
| **Remark 6:** The fact that we supposed $(a_{16}'')^{(2)}$ and $(b_{16}'')^{(2)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate | 199 |

condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{16})^{(2)}e^{(\widehat{M}_{16})^{(2)}t}$ and $(\widehat{Q}_{16})^{(2)}e^{(\widehat{M}_{16})^{(2)}t}$ respectively of $\mathbb{R}_+$.

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}, i = 16,17,18$ depend only on $T_{17}$ and respectively on $(G_{19})$(and not on t) and hypothesis can replaced by a usual Lipschitz condition.

**Remark 7:** There does not exist any t where $G_i(t) = 0$ and $T_i(t) = 0$ 200

it results

$$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(2)} - (a_i^{''})^{(2)}(T_{17}(s_{(16)}), s_{(16)})\} ds_{(16)}\right]} \geq 0$$

$$T_i(t) \geq T_i^0 e^{(-(b_i^{'})^{(2)}t)} > 0 \quad \text{for } t > 0$$

**Definition of** $\left((\widehat{M}_{16})^{(2)}\right)_1, \left((\widehat{M}_{16})^{(2)}\right)_2$ and $\left((\widehat{M}_{16})^{(2)}\right)_3$ : 201

**Remark 8:** if $G_{16}$ is bounded, the same property have also $G_{17}$ and $G_{18}$ . indeed if

$G_{16} < (\widehat{M}_{16})^{(2)}$ it follows $\frac{dG_{17}}{dt} \leq \left((\widehat{M}_{16})^{(2)}\right)_1 - (a_{17}^{'})^{(2)}G_{17}$ and by integrating

$$G_{17} \leq \left((\widehat{M}_{16})^{(2)}\right)_2 = G_{17}^0 + 2(a_{17})^{(2)}\left((\widehat{M}_{16})^{(2)}\right)_1/(a_{17}^{'})^{(2)}$$

In the same way , one can obtain

$$G_{18} \leq \left((\widehat{M}_{16})^{(2)}\right)_3 = G_{18}^0 + 2(a_{18})^{(2)}\left((\widehat{M}_{16})^{(2)}\right)_2/(a_{18}^{'})^{(2)}$$

If $G_{17}$ or $G_{18}$ is bounded, the same property follows for $G_{16}$ , $G_{18}$ and $G_{16}$ , $G_{17}$ respectively.

**Remark 9:** If $G_{16}$ is bounded, from below, the same property holds for $G_{17}$ and $G_{18}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{17}$ is bounded from below. 202

**Remark 10:** If $T_{16}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(2)}((G_{19})(t), t)) = (b_{17}^{'})^{(2)}$ then $T_{17} \to \infty$. 203

**Definition of** $(m)^{(2)}$ and $\varepsilon_2$ :

Indeed let $t_2$ be so that for $t > t_2$

$$(b_{17})^{(2)} - (b_i^{''})^{(2)}((G_{19})(t), t) < \varepsilon_2, T_{16}(t) > (m)^{(2)}$$

Then $\frac{dT_{17}}{dt} \geq (a_{17})^{(2)}(m)^{(2)} - \varepsilon_2 T_{17}$ which leads to 204

$T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{\varepsilon_2}\right)(1 - e^{-\varepsilon_2 t}) + T_{17}^0 e^{-\varepsilon_2 t}$ If we take t such that $e^{-\varepsilon_2 t} = \frac{1}{2}$ it results

$T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{2}\right), \quad t = \log\frac{2}{\varepsilon_2}$ By taking now $\varepsilon_2$ sufficiently small one sees that $T_{17}$ is unbounded. 205

The same property holds for $T_{18}$ if $\lim_{t\to\infty}(b_{18}^{''})^{(2)}\left((G_{19})(t), t\right) = (b_{18}^{'})^{(2)}$

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| We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |
| It is now sufficient to take $\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}$ , $\frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}} < 1$ and to choose<br>$(\hat{P}_{20})^{(3)}$ and $(\hat{Q}_{20})^{(3)}$ large to have | 207 |
| $\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}\left[(\hat{P}_{20})^{(3)} + \left((\hat{P}_{20})^{(3)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{20})^{(3)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{20})^{(3)}$ | 208 |
| $\frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}}\left[\left((\hat{Q}_{20})^{(3)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{20})^{(3)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{20})^{(3)}\right] \leq (\hat{Q}_{20})^{(3)}$ | 209 |
| In order that the operator $\mathcal{A}^{(3)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 210 |
| The operator $\mathcal{A}^{(3)}$ is a contraction with respect to the metric<br>$d\left(((G_{23})^{(1)}, (T_{23})^{(1)}), ((G_{23})^{(2)}, (T_{23})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\hat{M}_{20})^{(3)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\hat{M}_{20})^{(3)}t}\}$ | 211 |
| Indeed if we denote<br>**Definition of** $\widetilde{G_{23}}, \widetilde{T_{23}} :(\widetilde{(G_{23})}, \widetilde{(T_{23})}) = \mathcal{A}^{(3)}((G_{23}), (T_{23}))$ | 212 |
| It results<br>$\left\|\tilde{G}_{20}^{(1)} - \tilde{G}_i^{(2)}\right\| \leq \int_0^t (a_{20})^{(3)}\left\|G_{21}^{(1)} - G_{21}^{(2)}\right\|e^{-(\hat{M}_{20})^{(3)}s_{(20)}}e^{(\hat{M}_{20})^{(3)}s_{(20)}}\,ds_{(20)} +$<br>$\int_0^t\{(a'_{20})^{(3)}\left\|G_{20}^{(1)} - G_{20}^{(2)}\right\|e^{-(\hat{M}_{20})^{(3)}s_{(20)}}e^{-(\hat{M}_{20})^{(3)}s_{(20)}} +$<br>$(a''_{20})^{(3)}\left(T_{21}^{(1)}, s_{(20)}\right)\left\|G_{20}^{(1)} - G_{20}^{(2)}\right\|e^{-(\hat{M}_{20})^{(3)}s_{(20)}}e^{(\hat{M}_{20})^{(3)}s_{(20)}} +$<br>$G_{20}^{(2)}\left\|(a''_{20})^{(3)}\left(T_{21}^{(1)}, s_{(20)}\right) - (a''_{20})^{(3)}\left(T_{21}^{(2)}, s_{(20)}\right)\right\| \; e^{-(\hat{M}_{20})^{(3)}s_{(20)}}e^{(\hat{M}_{20})^{(3)}s_{(20)}}\}ds_{(20)}$<br>Where $s_{(20)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | 213 |
| $\left\|G_{23}{}^{(1)} - G_{23}{}^{(2)}\right\|e^{-(\hat{M}_{20})^{(3)}t} \leq$<br>$\frac{1}{(\hat{M}_{20})^{(3)}}\left((a_{20})^{(3)} + (a'_{20})^{(3)} + (\hat{A}_{20})^{(3)} + (\hat{P}_{20})^{(3)}(\hat{k}_{20})^{(3)}\right)d\left(((G_{23})^{(1)}, (T_{23})^{(1)};\ (G_{23})^{(2)}, (T_{23})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 214 |

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| **Remark 11:** The fact that we supposed $(a_{20}^{''})^{(3)}$ and $(b_{20}^{''})^{(3)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}$ $and$ $(\widehat{Q}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(3)}$ and $(b_i^{''})^{(3)}, i = 20,21,22$ depend only on $T_{21}$ and respectively on $(G_{23})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 215 |
| **Remark 12:** There does not exist any $t$ where $G_i(t) = 0\ and\ T_i(t) = 0$<br><br>it results<br><br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(3)} - (a_i^{''})^{(3)}(T_{21}(s_{(20)}), s_{(20)})\} ds_{(20)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{(-(b_i^{'})^{(3)}t)} > 0$ for t > 0 | 216 |
| **Definition of** $\left((\widehat{M}_{20})^{(3)}\right)_1, \left((\widehat{M}_{20})^{(3)}\right)_2\ and\ \left((\widehat{M}_{20})^{(3)}\right)_3$ :<br><br>**Remark 13:** if $G_{20}$ is bounded, the same property have also $G_{21}\ and\ G_{22}$ . indeed if<br><br>$G_{20} < (\widehat{M}_{20})^{(3)}$ it follows $\frac{dG_{21}}{dt} \leq \left((\widehat{M}_{20})^{(3)}\right)_1 - (a_{21}^{'})^{(3)}G_{21}$ and by integrating<br><br>$G_{21} \leq \left((\widehat{M}_{20})^{(3)}\right)_2 = G_{21}^0 + 2(a_{21})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_1/(a_{21}^{'})^{(3)}$<br><br>In the same way , one can obtain<br><br>$G_{22} \leq \left((\widehat{M}_{20})^{(3)}\right)_3 = G_{22}^0 + 2(a_{22})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_2/(a_{22}^{'})^{(3)}$<br><br>If $G_{21}\ or\ G_{22}$ is bounded, the same property follows for $G_{20}$ , $G_{22}$ and $G_{20}$ , $G_{21}$ respectively. | 217 |
| **Remark 14:** If $G_{20}$ $is$ bounded, from below, the same property holds for $G_{21}\ and\ G_{22}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{21}$ is bounded from below. | 218 |
| **Remark 15:** If $T_{20}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(3)}\left((G_{23})(t), t\right)) = (b_{21}^{'})^{(3)}$ then $T_{21} \to \infty$.<br><br>**Definition of** $(m)^{(3)}$ and $\varepsilon_3$ :<br><br>Indeed let $t_3$ be so that for $t > t_3$<br><br>$(b_{21})^{(3)} - (b_i^{''})^{(3)}\left((G_{23})(t), t\right) < \varepsilon_3, T_{20}(t) > (m)^{(3)}$ | 219 |
| Then $\frac{dT_{21}}{dt} \geq (a_{21})^{(3)}(m)^{(3)} - \varepsilon_3 T_{21}$ which leads to<br><br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{\varepsilon_3}\right)(1 - e^{-\varepsilon_3 t}) + T_{21}^0 e^{-\varepsilon_3 t}$ If we take t such that $e^{-\varepsilon_3 t} = \frac{1}{2}$ it results<br><br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{2}\right)$, $t = log\frac{2}{\varepsilon_3}$ By taking now $\varepsilon_3$ sufficiently small one sees that $T_{21}$ is unbounded. | 220 |

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| The same property holds for $T_{22}$ if $\lim_{t\to\infty}(b_{22}'')^{(3)}\left((G_{23})(t),t\right)=(b_{22}')^{(3)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |
| It is now sufficient to take $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}$ , $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}<1$ and to choose<br>$(\widehat{P}_{24})^{(4)}$ and $(\widehat{Q}_{24})^{(4)}$ large to have | 221 |
| $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[(\widehat{P}_{24})^{(4)}+\left((\hat{P}_{24})^{(4)}+G_j^0\right)e^{-\left(\frac{(\hat{P}_{24})^{(4)}+G_j^0}{G_j^0}\right)}\right]\leq(\hat{P}_{24})^{(4)}$ | 222 |
| $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[\left((\hat{Q}_{24})^{(4)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{24})^{(4)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{24})^{(4)}\right]\leq(\hat{Q}_{24})^{(4)}$ | 223 |
| In order that the operator $\mathcal{A}^{(4)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 224 |
| The operator $\mathcal{A}^{(4)}$ is a contraction with respect to the metric<br>$d\left(((G_{27})^{(1)},(T_{27})^{(1)}),((G_{27})^{(2)},(T_{27})^{(2)})\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{24})^{(4)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{24})^{(4)}t}\}$<br>Indeed if we denote<br>**Definition of** $(\widetilde{G_{27}}),(\widetilde{T_{27}})$ : $\left((\widetilde{G_{27}}),(\widetilde{T_{27}})\right)=\mathcal{A}^{(4)}((G_{27}),(T_{27}))$<br>It results<br>$\left\|\tilde{G}_{24}^{(1)}-\tilde{G}_i^{(2)}\right\|\leq\int_0^t(a_{24})^{(4)}\left\|G_{25}^{(1)}-G_{25}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}\,ds_{(24)}+$<br>$\int_0^t\{(a_{24}')^{(4)}\left\|G_{24}^{(1)}-G_{24}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}+$<br>$(a_{24}'')^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)\left\|G_{24}^{(1)}-G_{24}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}+$<br>$G_{24}^{(2)}\left\|(a_{24}'')^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)-(a_{24}'')^{(4)}\left(T_{25}^{(2)},s_{(24)}\right)\right\|\ e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}\}ds_{(24)}$<br>Where $s_{(24)}$ represents integrand that is integrated over the interval $[0,t]$<br>From the hypotheses on Equations it follows | 225 |
| $\left\|(G_{27})^{(1)}-(G_{27})^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}t}\leq$<br>$\frac{1}{(\widehat{M}_{24})^{(4)}}\left((a_{24})^{(4)}+(a_{24}')^{(4)}+(\widehat{A}_{24})^{(4)}+(\widehat{P}_{24})^{(4)}(\widehat{k}_{24})^{(4)}\right)d\left(((G_{27})^{(1)},(T_{27})^{(1)};\ (G_{27})^{(2)},(T_{27})^{(2)})\right)$ | 226 |

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| And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | |
| **Remark 16:** The fact that we supposed $(a_{24}^{''})^{(4)}$ and $(b_{24}^{''})^{(4)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}$ *and* $(\widehat{Q}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(4)}$ and $(b_i^{''})^{(4)}, i = 24,25,26$ depend only on $T_{25}$ and respectively on $(G_{27})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 227 |
| **Remark 17:** There does not exist any $t$ where $G_i\,(t) = 0\ and\ T_i\,(t) = 0$<br>it results<br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t\{(a_i^{'})^{(4)}-(a_i^{''})^{(4)}(T_{25}(s_{(24)}),s_{(24)})\}ds_{(24)}\right]} \geq 0$<br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(4)}t\right)} > 0$ for t > 0 | 228 |
| **Definition of** $\left((\widehat{M}_{24})^{(4)}\right)_1, \left((\widehat{M}_{24})^{(4)}\right)_2\ and\ \left((\widehat{M}_{24})^{(4)}\right)_3$ :<br>**Remark 18:** if $G_{24}$ is bounded, the same property have also $G_{25}\ and\ G_{26}$ . indeed if<br>$G_{24} < (\widehat{M}_{24})^{(4)}$ it follows $\frac{dG_{25}}{dt} \leq \left((\widehat{M}_{24})^{(4)}\right)_1 - (a_{25}^{'})^{(4)}G_{25}$ and by integrating<br>$G_{25} \leq \left((\widehat{M}_{24})^{(4)}\right)_2 = G_{25}^0 + 2(a_{25})^{(4)}\left((\widehat{M}_{24})^{(4)}\right)_1/(a_{25}^{'})^{(4)}$<br>In the same way , one can obtain<br>$G_{26} \leq \left((\widehat{M}_{24})^{(4)}\right)_3 = G_{26}^0 + 2(a_{26})^{(4)}\left((\widehat{M}_{24})^{(4)}\right)_2/(a_{26}^{'})^{(4)}$<br>If $G_{25}\ or\ G_{26}$ is bounded, the same property follows for $G_{24}$ , $G_{26}$ and $G_{24}$ , $G_{25}$ respectively. | 229 |
| **Remark 19:** If $G_{24}$ *is* bounded, from below, the same property holds for $G_{25}\ and\ G_{26}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{25}$ is bounded from below. | 230 |
| **Remark 20:** If $T_{24}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(4)}\,((G_{27})(t),t)) = (b_{25}^{'})^{(4)}$ then $T_{25} \to \infty$.<br>**Definition of** $(m)^{(4)}$ and $\varepsilon_4$ :<br>Indeed let $t_4$ be so that for $t > t_4$<br>$(b_{25})^{(4)} - (b_i^{''})^{(4)}((G_{27})(t),t) < \varepsilon_4, T_{24}\,(t) > (m)^{(4)}$ | 231 |
| Then $\frac{dT_{25}}{dt} \geq (a_{25})^{(4)}(m)^{(4)} - \varepsilon_4 T_{25}$ which leads to<br>$T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{\varepsilon_4}\right)(1 - e^{-\varepsilon_4 t}) + T_{25}^0 e^{-\varepsilon_4 t}$ If we take t such that $e^{-\varepsilon_4 t} = \frac{1}{2}$ it results | 232 |

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| $T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{2}\right), \quad t = log\frac{2}{\varepsilon_4}$ By taking now $\varepsilon_4$ sufficiently small one sees that $T_{25}$ is unbounded. The same property holds for $T_{26}$ if $\lim_{t\to\infty}(b_{26}^{''})^{(4)}\left((G_{27})(t),t\right) = (b_{26}^{'})^{(4)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 42<br><br>Analogous inequalities hold also for $G_{29}, G_{30}, T_{28}, T_{29}, T_{30}$ | |
| It is now sufficient to take $\frac{(a_i)^{(5)}}{(\hat{M}_{28})^{(5)}}, \frac{(b_i)^{(5)}}{(\hat{M}_{28})^{(5)}} < 1$ and to choose<br><br>$(\hat{P}_{28})^{(5)}$ and $(\hat{Q}_{28})^{(5)}$ large to have | 233 |
| $\frac{(a_i)^{(5)}}{(\hat{M}_{28})^{(5)}}\left[(\hat{P}_{28})^{(5)} + \left((\hat{P}_{28})^{(5)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{28})^{(5)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{28})^{(5)}$ | 234 |
| $\frac{(b_i)^{(5)}}{(\hat{M}_{28})^{(5)}}\left[\left((\hat{Q}_{28})^{(5)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{28})^{(5)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{28})^{(5)}\right] \leq (\hat{Q}_{28})^{(5)}$ | 235 |
| In order that the operator $\mathcal{A}^{(5)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(5)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{31})^{(1)},(T_{31})^{(1)}),((G_{31})^{(2)},(T_{31})^{(2)})\right) =$<br><br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t) - G_i^{(2)}(t)\right|e^{-(\hat{M}_{28})^{(5)}t}, \max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t) - T_i^{(2)}(t)\right|e^{-(\hat{M}_{28})^{(5)}t}\}$<br><br>Indeed if we denote<br><br>**Definition of** $(\widetilde{G_{31}}),(\widetilde{T_{31}})$ : $\left((\widetilde{G_{31}}),(\widetilde{T_{31}})\right) = \mathcal{A}^{(5)}\left((G_{31}),(T_{31})\right)$<br><br>It results<br><br>$\left|\tilde{G}_{28}^{(1)} - \tilde{G}_i^{(2)}\right| \leq \int_0^t (a_{28})^{(5)}\left|G_{29}^{(1)} - G_{29}^{(2)}\right|e^{-(\hat{M}_{28})^{(5)}s_{(28)}}e^{(\hat{M}_{28})^{(5)}s_{(28)}}\,ds_{(28)} +$<br><br>$\int_0^t\{(a_{28}^{'})^{(5)}\left|G_{28}^{(1)} - G_{28}^{(2)}\right|e^{-(\hat{M}_{28})^{(5)}s_{(28)}}e^{-(\hat{M}_{28})^{(5)}s_{(28)}} +$<br><br>$(a_{28}^{''})^{(5)}\left(T_{29}^{(1)},s_{(28)}\right)\left|G_{28}^{(1)} - G_{28}^{(2)}\right|e^{-(\hat{M}_{28})^{(5)}s_{(28)}}e^{(\hat{M}_{28})^{(5)}s_{(28)}} +$<br><br>$G_{28}^{(2)}\left|(a_{28}^{''})^{(5)}\left(T_{29}^{(1)},s_{(28)}\right) - (a_{28}^{''})^{(5)}\left(T_{29}^{(2)},s_{(28)}\right)\right|\ e^{-(\hat{M}_{28})^{(5)}s_{(28)}}e^{(\hat{M}_{28})^{(5)}s_{(28)}}\}ds_{(28)}$<br><br>Where $s_{(28)}$ represents integrand that is integrated over the interval [0, t] | 236 |

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| From the hypotheses on it follows | |
| $\left\|(G_{31})^{(1)} - (G_{31})^{(2)}\right\| e^{-(\widehat{M}_{28})^{(5)}t} \leq$ $\frac{1}{(\widehat{M}_{28})^{(5)}}\left((a_{28})^{(5)} + (a'_{28})^{(5)} + (\widehat{A}_{28})^{(5)} + (\widehat{P}_{28})^{(5)}(\widehat{k}_{28})^{(5)}\right) d\left(((G_{31})^{(1)}, (T_{31})^{(1)};\ (G_{31})^{(2)}, (T_{31})^{(2)})\right)$ And analogous inequalities for $G_i$ and $T_i$. Taking into account the hypothesis the result follows | 237 |
| **Remark 21:** The fact that we supposed $(a''_{28})^{(5)}$ and $(b''_{28})^{(5)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{28})^{(5)} e^{(\widehat{M}_{28})^{(5)}t}$ and $(\widehat{Q}_{28})^{(5)} e^{(\widehat{M}_{28})^{(5)}t}$ respectively of $\mathbb{R}_+$. If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a''_i)^{(5)}$ and $(b''_i)^{(5)}, i = 28,29,30$ depend only on $T_{29}$ and respectively on $(G_{31})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 238 |
| **Remark 22:** There does not exist any $t$ where $G_i(t) = 0$ and $T_i(t) = 0$ it results $G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a'_i)^{(5)} - (a''_i)^{(5)}(T_{29}(s_{(28)}), s_{(28)})\} ds_{(28)}\right]} \geq 0$ $T_i(t) \geq T_i^0 e^{\left(-(b'_i)^{(5)}t\right)} > 0$ for t > 0 | 239 |
| **Definition of** $\left((\widehat{M}_{28})^{(5)}\right)_1, \left((\widehat{M}_{28})^{(5)}\right)_2$ and $\left((\widehat{M}_{28})^{(5)}\right)_3$ : **Remark 23:** if $G_{28}$ is bounded, the same property have also $G_{29}$ and $G_{30}$ . indeed if $G_{28} < (\widehat{M}_{28})^{(5)}$ it follows $\frac{dG_{29}}{dt} \leq \left((\widehat{M}_{28})^{(5)}\right)_1 - (a'_{29})^{(5)} G_{29}$ and by integrating $G_{29} \leq \left((\widehat{M}_{28})^{(5)}\right)_2 = G_{29}^0 + 2(a_{29})^{(5)}\left((\widehat{M}_{28})^{(5)}\right)_1 / (a'_{29})^{(5)}$ In the same way , one can obtain $G_{30} \leq \left((\widehat{M}_{28})^{(5)}\right)_3 = G_{30}^0 + 2(a_{30})^{(5)}\left((\widehat{M}_{28})^{(5)}\right)_2 / (a'_{30})^{(5)}$ If $G_{29}$ or $G_{30}$ is bounded, the same property follows for $G_{28}$ , $G_{30}$ and $G_{28}$ , $G_{29}$ respectively. | 240 |
| **Remark 24:** If $G_{28}$ is bounded, from below, the same property holds for $G_{29}$ and $G_{30}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{29}$ is bounded from below. | 241 |

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| **Remark 25:** If $T_{28}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(5)}((G_{31})(t),t)) = (b_{29}^{'})^{(5)}$ then $T_{29} \to \infty$.<br><br>**Definition of** $(m)^{(5)}$ and $\varepsilon_5$ :<br><br>Indeed let $t_5$ be so that for $t > t_5$<br><br>$(b_{29})^{(5)} - (b_i^{''})^{(5)}((G_{31})(t),t) < \varepsilon_5, T_{28}(t) > (m)^{(5)}$ | 242 |
| Then $\frac{dT_{29}}{dt} \geq (a_{29})^{(5)}(m)^{(5)} - \varepsilon_5 T_{29}$ which leads to<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{\varepsilon_5}\right)(1-e^{-\varepsilon_5 t}) + T_{29}^0 e^{-\varepsilon_5 t}$ If we take t such that $e^{-\varepsilon_5 t} = \frac{1}{2}$ it results<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{2}\right), \quad t = log\frac{2}{\varepsilon_5}$ By taking now $\varepsilon_5$ sufficiently small one sees that $T_{29}$ is unbounded.<br>The same property holds for $T_{30}$ if $\lim_{t\to\infty}(b_{30}^{''})^{(5)}\left((G_{31})(t),t\right) = (b_{30}^{'})^{(5)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations<br><br>Analogous inequalities hold also for $G_{33}, G_{34}, T_{32}, T_{33}, T_{34}$ | 243 |
| It is now sufficient to take $\frac{(a_i)^{(6)}}{(\hat{M}_{32})^{(6)}}, \frac{(b_i)^{(6)}}{(\hat{M}_{32})^{(6)}} < 1$ and to choose<br><br>$(\hat{P}_{32})^{(6)}$ and $(\hat{Q}_{32})^{(6)}$ large to have | 244 |
| $\frac{(a_i)^{(6)}}{(\hat{M}_{32})^{(6)}}\left[(\hat{P}_{32})^{(6)} + \left((\hat{P}_{32})^{(6)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{32})^{(6)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{32})^{(6)}$ | 245 |
| $\frac{(b_i)^{(6)}}{(\hat{M}_{32})^{(6)}}\left[\left((\hat{Q}_{32})^{(6)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{32})^{(6)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{32})^{(6)}\right] \leq (\hat{Q}_{32})^{(6)}$ | 246 |
| In order that the operator $\mathcal{A}^{(6)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(6)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{35})^{(1)},(T_{35})^{(1)}),((G_{35})^{(2)},(T_{35})^{(2)})\right) =$<br><br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t) - G_i^{(2)}(t)\right|e^{-(\hat{M}_{32})^{(6)}t}, \max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t) - T_i^{(2)}(t)\right|e^{-(\hat{M}_{32})^{(6)}t}\}$<br><br>Indeed if we denote<br><br>**Definition of** $\widetilde{(G_{35})}, \widetilde{(T_{35})}$ : $\left(\widetilde{(G_{35})}, \widetilde{(T_{35})}\right) = \mathcal{A}^{(6)}\left((G_{35}),(T_{35})\right)$<br><br>It results | 247 |

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| $\left\|\tilde{G}_{32}^{(1)}-\tilde{G}_{i}^{(2)}\right\| \leq \int_0^t (a_{32})^{(6)} \left\|G_{33}^{(1)}-G_{33}^{(2)}\right\| e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} e^{(\widehat{M}_{32})^{(6)}s_{(32)}}\, ds_{(32)} +$<br>$\int_0^t \{(a'_{32})^{(6)} \left\|G_{32}^{(1)}-G_{32}^{(2)}\right\| e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} +$<br>$(a''_{32})^{(6)}\left(T_{33}^{(1)}, s_{(32)}\right)\left\|G_{32}^{(1)}-G_{32}^{(2)}\right\| e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} e^{(\widehat{M}_{32})^{(6)}s_{(32)}} +$<br>$G_{32}^{(2)} \left\|(a''_{32})^{(6)}\left(T_{33}^{(1)}, s_{(32)}\right) - (a''_{32})^{(6)}\left(T_{33}^{(2)}, s_{(32)}\right)\right\| \; e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} e^{(\widehat{M}_{32})^{(6)}s_{(32)}}\} ds_{(32)}$<br>Where $s_{(32)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | |
| $\left\|(G_{35})^{(1)}-(G_{35})^{(2)}\right\| e^{-(\widehat{M}_{32})^{(6)}t} \leq$<br>$\frac{1}{(\widehat{M}_{32})^{(6)}}\left((a_{32})^{(6)} + (a'_{32})^{(6)} + (\widehat{A}_{32})^{(6)} + (\widehat{P}_{32})^{(6)}(\widehat{k}_{32})^{(6)}\right) d\left(((G_{35})^{(1)}, (T_{35})^{(1)};\ (G_{35})^{(2)}, (T_{35})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 248 |
| **Remark 26:** The fact that we supposed $(a''_{32})^{(6)}$ and $(b''_{32})^{(6)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{32})^{(6)} e^{(\widehat{M}_{32})^{(6)}t}$ *and* $(\widehat{Q}_{32})^{(6)} e^{(\widehat{M}_{32})^{(6)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i'')^{(6)}$ and $(b_i'')^{(6)}, i = 32,33,34$ depend only on $T_{33}$ and respectively on $(G_{35})$(*and not on* $t$) and hypothesis can replaced by a usual Lipschitz condition. | 249 |
| **Remark 27:** There does not exist any $t$ where $G_i(t) = 0$ *and* $T_i(t) = 0$<br>it results<br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i')^{(6)} - (a_i'')^{(6)}(T_{33}(s_{(32)}), s_{(32)})\} ds_{(32)}\right]} \geq 0$<br>$T_i(t) \geq T_i^0 e^{\left(-(b_i')^{(6)}t\right)} > 0$ for t > 0 | 250 |
| **Definition of** $\left((\widehat{M}_{32})^{(6)}\right)_1, \left((\widehat{M}_{32})^{(6)}\right)_2$ *and* $\left((\widehat{M}_{32})^{(6)}\right)_3$ :<br>**Remark 28:** if $G_{32}$ is bounded, the same property have also $G_{33}$ *and* $G_{34}$ . indeed if<br>$G_{32} < (\widehat{M}_{32})^{(6)}$ it follows $\frac{dG_{33}}{dt} \leq \left((\widehat{M}_{32})^{(6)}\right)_1 - (a'_{33})^{(6)} G_{33}$ and by integrating<br>$G_{33} \leq \left((\widehat{M}_{32})^{(6)}\right)_2 = G_{33}^0 + 2(a_{33})^{(6)}\left((\widehat{M}_{32})^{(6)}\right)_1 / (a'_{33})^{(6)}$<br>In the same way , one can obtain<br>$G_{34} \leq \left((\widehat{M}_{32})^{(6)}\right)_3 = G_{34}^0 + 2(a_{34})^{(6)}\left((\widehat{M}_{32})^{(6)}\right)_2 / (a'_{34})^{(6)}$ | 251 |

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| If $G_{33}$ *or* $G_{34}$ is bounded, the same property follows for $G_{32}$ , $G_{34}$ and $G_{32}$ , $G_{33}$ respectively. | |
| **Remark 29:** If $G_{32}$ *is* bounded, from below, the same property holds for $G_{33}$ *and* $G_{34}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{33}$ is bounded from below. | 252 |
| **Remark 30:** If $T_{32}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(6)}((G_{35})(t),t)) = (b_{33}^{'})^{(6)}$ then $T_{33} \to \infty$.<br><br>**Definition of** $(m)^{(6)}$ and $\varepsilon_6$ :<br><br>Indeed let $t_6$ be so that for $t > t_6$<br><br>$(b_{33})^{(6)} - (b_i^{''})^{(6)}\big((G_{35})(t),t\big) < \varepsilon_6, T_{32}(t) > (m)^{(6)}$ | 253 |
| Then $\frac{dT_{33}}{dt} \geq (a_{33})^{(6)}(m)^{(6)} - \varepsilon_6 T_{33}$ which leads to<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{\varepsilon_6}\right)(1 - e^{-\varepsilon_6 t}) + T_{33}^0 e^{-\varepsilon_6 t}$ If we take t such that $e^{-\varepsilon_6 t} = \frac{1}{2}$ it results<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{2}\right)$, $t = log\frac{2}{\varepsilon_6}$ By taking now $\varepsilon_6$ sufficiently small one sees that $T_{33}$ is unbounded.<br>The same property holds for $T_{34}$ if $\lim_{t\to\infty}(b_{34}^{''})^{(6)}\left((G_{35})(t),t(t),t\right) = (b_{34}^{'})^{(6)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 254 |
| Analogous inequalities hold also for $G_{37}$ , $G_{38}, T_{36}, T_{37}, T_{38}$<br><br>It is now sufficient to take $\frac{(a_i)^{(7)}}{(\hat{M}_{36})^{(7)}}$ , $\frac{(b_i)^{(7)}}{(\hat{M}_{36})^{(7)}} < 1$ and to choose<br>$(\hat{P}_{36})^{(7)}$ and $(\hat{Q}_{36})^{(7)}$ large to have | 255 |
| $\frac{(a_i)^{(7)}}{(\hat{M}_{36})^{(7)}}\left[(\hat{P}_{36})^{(7)} + \left((\hat{P}_{36})^{(7)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{36})^{(7)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{36})^{(7)}$ | 256 |
| $\frac{(b_i)^{(7)}}{(\hat{M}_{36})^{(7)}}\left[\left((\hat{Q}_{36})^{(7)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{36})^{(7)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{36})^{(7)}\right] \leq (\hat{Q}_{36})^{(7)}$ | 257 |
| In order that the operator $\mathcal{A}^{(7)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(7)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{39})^{(1)},(T_{39})^{(1)}),((G_{39})^{(2)},(T_{39})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\hat{M}_{36})^{(7)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\hat{M}_{36})^{(7)}t}\}$<br><br>Indeed if we denote | 258 |

**Definition of** $\widetilde{(G_{39})}, \widetilde{(T_{39})} : \left( \widetilde{(G_{39})}, \widetilde{(T_{39})} \right) = \mathcal{A}^{(7)}((G_{39}), (T_{39}))$
It results

$\left|\tilde{G}_{36}^{(1)} - \tilde{G}_{i}^{(2)}\right| \leq \int_0^t (a_{36})^{(7)} \left|G_{37}^{(1)} - G_{37}^{(2)}\right| e^{-(\widehat{M}_{36})^{(7)} s_{(36)}} e^{(\widehat{M}_{36})^{(7)} s_{(36)}}\, ds_{(36)} +$
$\int_0^t \{(a_{36}^{'})^{(7)} \left|G_{36}^{(1)} - G_{36}^{(2)}\right| e^{-(\widehat{M}_{36})^{(7)} s_{(36)}} e^{-(\widehat{M}_{36})^{(7)} s_{(36)}} +$
$(a_{36}^{''})^{(7)}\left(T_{37}^{(1)}, s_{(36)}\right) \left|G_{36}^{(1)} - G_{36}^{(2)}\right| e^{-(\widehat{M}_{36})^{(7)} s_{(36)}} e^{(\widehat{M}_{36})^{(7)} s_{(36)}} +$
$G_{36}^{(2)} \left|(a_{36}^{''})^{(7)}\left(T_{37}^{(1)}, s_{(36)}\right) - (a_{36}^{''})^{(7)}\left(T_{37}^{(2)}, s_{(36)}\right)\right| \; e^{-(\widehat{M}_{36})^{(7)} s_{(36)}} e^{(\widehat{M}_{36})^{(7)} s_{(36)}} \} ds_{(36)}$

Where $s_{(36)}$ represents integrand that is integrated over the interval $[0, t]$

From the hypotheses on it follows

$\left|(G_{39})^{(1)} - (G_{39})^{(2)}\right| e^{-(\widehat{M}_{36})^{(7)} t} \leq$
$\frac{1}{(\widehat{M}_{36})^{(7)}} \left((a_{36})^{(7)} + (a_{36}^{'})^{(7)} + (\widehat{A}_{36})^{(7)} + (\widehat{P}_{36})^{(7)} (\widehat{k}_{36})^{(7)}\right) d\left(((G_{39})^{(1)}, (T_{39})^{(1)};\ (G_{39})^{(2)}, (T_{39})^{(2)})\right)$

And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows

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**Remark 31:** The fact that we supposed $(a_{36}^{''})^{(7)}$ and $(b_{36}^{''})^{(7)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{36})^{(7)} e^{(\widehat{M}_{36})^{(7)} t}$ *and* $(\widehat{Q}_{36})^{(7)} e^{(\widehat{M}_{36})^{(7)} t}$ respectively of $\mathbb{R}_+$.

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(7)}$ and $(b_i^{''})^{(7)}, i = 36,37,38$ depend only on $T_{37}$ and respectively on $(G_{39})$(*and not on* $t$) and hypothesis can replaced by a usual Lipschitz condition.

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**Remark 32:** There does not exist any $t$ where $G_i(t) = 0$ *and* $T_i(t) = 0$

it results

$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(7)} - (a_i^{''})^{(7)} (T_{37}(s_{(36)}), s_{(36)})\} ds_{(36)}\right]} \geq 0$

$T_i(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(7)} t\right)} > 0$ for t > 0

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**Definition of** $\left((\widehat{M}_{36})^{(7)}\right)_1, \left((\widehat{M}_{36})^{(7)}\right)_2$ *and* $\left((\widehat{M}_{36})^{(7)}\right)_3$ :

**Remark 33:** if $G_{36}$ is bounded, the same property have also $G_{37}$ *and* $G_{38}$ . indeed if

$G_{36} < (\widehat{M}_{36})^{(7)}$ it follows $\frac{dG_{37}}{dt} \leq \left((\widehat{M}_{36})^{(7)}\right)_1 - (a_{37}^{'})^{(7)} G_{37}$ and by integrating

$G_{37} \leq \left((\widehat{M}_{36})^{(7)}\right)_2 = G_{37}^0 + 2(a_{37})^{(7)} \left((\widehat{M}_{36})^{(7)}\right)_1 / (a_{37}^{'})^{(7)}$

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| In the same way , one can obtain<br>$G_{38} \leq \left((\widehat{M}_{36})^{(7)}\right)_3 = G_{38}^0 + 2(a_{38})^{(7)}\left((\widehat{M}_{36})^{(7)}\right)_2/(a_{38}')^{(7)}$<br>If $G_{37}$ *or* $G_{38}$ is bounded, the same property follows for $G_{36}$ , $G_{38}$ and $G_{36}$ , $G_{37}$ respectively. | |
| **Remark 34:** If $G_{36}$ *is* bounded, from below, the same property holds for $G_{37}$ *and* $G_{38}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{37}$ is bounded from below. | 263 |
| **Remark 35:** If $T_{36}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(7)}((G_{39})(t),t)) = (b_{37}')^{(7)}$ then $T_{37} \to \infty$.<br>**Definition of** $(m)^{(7)}$ and $\varepsilon_7$ :<br>Indeed let $t_7$ be so that for $t > t_7$<br>$(b_{37})^{(7)} - (b_i'')^{(7)}((G_{39})(t),t) < \varepsilon_7, T_{36}(t) > (m)^{(7)}$ | 264 |
| Then $\frac{dT_{37}}{dt} \geq (a_{37})^{(7)}(m)^{(7)} - \varepsilon_7 T_{37}$ which leads to<br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{\varepsilon_7}\right)(1 - e^{-\varepsilon_7 t}) + T_{37}^0 e^{-\varepsilon_7 t}$ If we take t such that $e^{-\varepsilon_7 t} = \frac{1}{2}$ it results<br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{2}\right)$, $t = log\frac{2}{\varepsilon_7}$ By taking now $\varepsilon_7$ sufficiently small one sees that $T_{37}$ is unbounded.<br>The same property holds for $T_{38}$ if $\lim_{t\to\infty}(b_{38}'')^{(7)}\left((G_{39})(t),t\right) = (b_{38}')^{(7)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 265 |
| It is now sufficient to take $\frac{(a_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}$ , $\frac{(b_i)^{(8)}}{(\widehat{M}_{40})^{(8)}} < 1$ and to choose<br>$(\widehat{P}_{40})^{(8)}$ and $(\widehat{Q}_{40})^{(8)}$ large to have | 266 |
| $\frac{(a_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}\left[(\widehat{P}_{40})^{(8)} + \left((\widehat{P}_{40})^{(8)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{40})^{(8)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{40})^{(8)}$ | 267 |
| $\frac{(b_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}\left[\left((\widehat{Q}_{40})^{(8)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{40})^{(8)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{40})^{(8)}\right] \leq (\widehat{Q}_{40})^{(8)}$ | 268 |
| In order that the operator $\mathcal{A}^{(8)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(8)}$ is a contraction with respect to the metric | |

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| $d\left(((G_{43})^{(1)},(T_{43})^{(1)}),((G_{43})^{(2)},(T_{43})^{(2)})\right)=$ $\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{40})^{(8)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{40})^{(8)}t}\}$ | 269 |
| Indeed if we denote<br><br>**Definition of** $\widetilde{(G_{43})},\widetilde{(T_{43})}$ : $\left(\widetilde{(G_{43})},\widetilde{(T_{43})}\right)=\mathcal{A}^{(8)}((G_{43}),(T_{43}))$ | 270 |
| It results<br><br>$\left\|\tilde{G}_{40}^{(1)}-\tilde{G}_i^{(2)}\right\|\le\int_0^t(a_{40})^{(8)}\left\|G_{41}^{(1)}-G_{41}^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{(\widehat{M}_{40})^{(8)}s_{(40)}}\,ds_{(40)}+$<br>$\int_0^t\{(a_{40}^{'})^{(8)}\left\|G_{40}^{(1)}-G_{40}^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}+$<br>$(a_{40}^{''})^{(8)}\left(T_{41}^{(1)},s_{(40)}\right)\left\|G_{40}^{(1)}-G_{40}^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{(\widehat{M}_{40})^{(8)}s_{(40)}}+$<br>$G_{40}^{(2)}\left\\|(a_{40}^{''})^{(8)}\left(T_{41}^{(1)},s_{(40)}\right)-(a_{40}^{''})^{(8)}\left(T_{41}^{(2)},s_{(40)}\right)\right\\|\ e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{(\widehat{M}_{40})^{(8)}s_{(40)}}\}ds_{(40)}$ | 271 |
| Where $s_{(40)}$ represents integrand that is integrated over the interval $[0,\mathrm{t}]$<br><br>From the hypotheses it follows | 272 |
| $\left\|(G_{43})^{(1)}-(G_{43})^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}t}\le$<br>$\frac{1}{(\widehat{M}_{40})^{(8)}}\left((a_{40})^{(8)}+(a_{40}^{'})^{(8)}+(\widehat{A}_{40})^{(8)}+(\widehat{P}_{40})^{(8)}(\widehat{k}_{40})^{(8)}\right)d\left(((G_{43})^{(1)},(T_{43})^{(1)};\ (G_{43})^{(2)},(T_{43})^{(2)})\right)$<br><br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 273 |
| **Remark 36:** The fact that we supposed $(a_{40}^{''})^{(8)}$ and $(b_{40}^{''})^{(8)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{40})^{(8)}e^{(\widehat{M}_{40})^{(8)}t}$ *and* $(\widehat{Q}_{40})^{(8)}e^{(\widehat{M}_{40})^{(8)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(8)}$ and $(b_i^{''})^{(8)}, i=40,41,42$ depend only on $\mathrm{T}_{41}$ and respectively on $(G_{43})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 274 |
| **Remark 37** There does not exist any $t$ where $G_i\,(t)=0$ *and* $T_i\,(t)=0$<br><br>it results<br><br>$G_i\,(t)\ge G_i^0e^{\left[-\int_0^t\{(a_i^{'})^{(8)}-(a_i^{''})^{(8)}(T_{41}(s_{(40)}),s_{(40)})\}ds_{(40)}\right]}\ge 0$<br><br>$T_i\,(t)\ge T_i^0e^{\left(-(b_i^{'})^{(8)}t\right)}>0$ for t > 0 | 275 |
| **Definition of** $\left((\widehat{M}_{40})^{(8)}\right)_1,\left((\widehat{M}_{40})^{(8)}\right)_2$ *and* $\left((\widehat{M}_{40})^{(8)}\right)_3$ : | 276 |

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| **Remark 38:** if $G_{40}$ is bounded, the same property have also $G_{41}\ and\ G_{42}$ . indeed if<br><br>$G_{40} < (\widehat{M}_{40})^{(8)}$ it follows $\frac{dG_{41}}{dt} \leq \left((\widehat{M}_{40})^{(8)}\right)_1 - (a'_{41})^{(8)}G_{41}$ and by integrating<br><br>$G_{41} \leq \left((\widehat{M}_{40})^{(8)}\right)_2 = G_{41}^0 + 2(a_{41})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_1/(a'_{41})^{(8)}$<br><br>In the same way , one can obtain<br><br>$G_{42} \leq \left((\widehat{M}_{40})^{(8)}\right)_3 = G_{42}^0 + 2(a_{42})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_2/(a'_{42})^{(8)}$<br><br>If $G_{41}\ or\ G_{42}$ is bounded, the same property follows for $G_{40}$ , $G_{42}$ and $G_{40}$ , $G_{41}$ respectively. | |
| **Remark 39:** If $G_{40}$ *is* bounded, from below, the same property holds for $G_{41}\ and\ G_{42}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{41}$ is bounded from below. | 277 |
| **Remark 40:** If $T_{40}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(8)}((G_{43})(t),t)) = (b'_{41})^{(8)}$ then $T_{41} \to \infty$.<br><br>**Definition of** $(m)^{(8)}$ and $\varepsilon_8$ :<br><br>Indeed let $t_8$ be so that for $t > t_8$<br><br>$(b_{41})^{(8)} - (b_i^{''})^{(8)}\left((G_{43})(t),t\right) < \varepsilon_8, T_{40}(t) > (m)^{(8)}$ | 278 |
| Then $\frac{dT_{41}}{dt} \geq (a_{41})^{(8)}(m)^{(8)} - \varepsilon_8 T_{41}$ which leads to<br><br>$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{\varepsilon_8}\right)(1 - e^{-\varepsilon_8 t}) + T_{41}^0 e^{-\varepsilon_8 t}$ If we take t such that $e^{-\varepsilon_8 t} = \frac{1}{2}$ it results<br><br>$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{2}\right),\ t = log\frac{2}{\varepsilon_8}$ By taking now $\varepsilon_8$ sufficiently small one sees that $T_{41}$ is unbounded.<br>The same property holds for $T_{42}$ if $\lim_{t\to\infty}(b_{42}^{''})^{(8)}\left((G_{43})(t),t(t),t\right) = (b'_{42})^{(8)}$ | 279 |
| It is now sufficient to take $\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}$ , $\frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}} < 1$ and to choose $(\widehat{P}_{44})^{(9)}$ and $(\widehat{Q}_{44})^{(9)}$ large to have | 279 A |
| $\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[(\widehat{P}_{44})^{(9)} + \left((\widehat{P}_{44})^{(9)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{44})^{(9)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{44})^{(9)}$ | |
| $\frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[\left((\widehat{Q}_{44})^{(9)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{44})^{(9)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{44})^{(9)}\right] \leq (\widehat{Q}_{44})^{(9)}$ | |

In order that the operator $\mathcal{A}^{(9)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying 39,35,36 into itself

The operator $\mathcal{A}^{(9)}$ is a contraction with respect to the metric

$$d\left(((G_{47})^{(1)},(T_{47})^{(1)}),((G_{47})^{(2)},(T_{47})^{(2)})\right)=$$
$$sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t)-G_i^{(2)}(t)\right|e^{-(\widehat{M}_{44})^{(9)}t},\max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t)-T_i^{(2)}(t)\right|e^{-(\widehat{M}_{44})^{(9)}t}\}$$

Indeed if we denote

**Definition of** $\widetilde{(G_{47})},\widetilde{(T_{47})}:\left(\widetilde{(G_{47})},\widetilde{(T_{47})}\right)=\mathcal{A}^{(9)}\left((G_{47}),(T_{47})\right)$
It results

$$\left|\tilde{G}_{44}^{(1)}-\tilde{G}_i^{(2)}\right|\le\int_0^t(a_{44})^{(9)}\left|G_{45}^{(1)}-G_{45}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\,ds_{(44)}+$$
$$\int_0^t\{(a_{44}^{'})^{(9)}\left|G_{44}^{(1)}-G_{44}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}+$$
$$(a_{44}^{''})^{(9)}\left(T_{45}^{(1)},s_{(44)}\right)\left|G_{44}^{(1)}-G_{44}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}+$$
$$G_{44}^{(2)}\left|(a_{44}^{''})^{(9)}\left(T_{45}^{(1)},s_{(44)}\right)-(a_{44}^{''})^{(9)}\left(T_{45}^{(2)},s_{(44)}\right)\right|\ e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\}ds_{(44)}$$

Where $s_{(44)}$ represents integrand that is integrated over the interval $[0,\mathrm{t}]$

From the hypotheses on 45,46,47,28 and 29 it follows

$$\left|(G_{47})^{(1)}-G^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}t}\le$$
$$\frac{1}{(\widehat{M}_{44})^{(9)}}\left((a_{44})^{(9)}+(a_{44}^{'})^{(9)}+(\widehat{A}_{44})^{(9)}+(\widehat{P}_{44})^{(9)}(\widehat{k}_{44})^{(9)}\right)d\left(((G_{47})^{(1)},(T_{47})^{(1)};\ (G_{47})^{(2)},(T_{47})^{(2)})\right)$$

And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis (39,35,36) the result follows

**Remark 41:** The fact that we supposed $(a_{44}^{''})^{(9)}$ and $(b_{44}^{''})^{(9)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{44})^{(9)}e^{(\widehat{M}_{44})^{(9)}t}$ *and* $(\widehat{Q}_{44})^{(9)}e^{(\widehat{M}_{44})^{(9)}t}$ respectively of $\mathbb{R}_+$.

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(9)}$ and $(b_i^{''})^{(9)}, i=44{,}45{,}46$ depend only on $\mathrm{T}_{45}$ and respectively on $(G_{47})$(*and not on* $t$) and hypothesis can replaced by a usual Lipschitz condition.

**Remark 42:** There does not exist any $t$ where $G_i(t)=0$ *and* $T_i(t)=0$

From 99 to 44 it results

$$G_i(t)\ge G_i^0e^{\left[-\int_0^t\{(a_i^{'})^{(9)}-(a_i^{''})^{(9)}(T_{45}(s_{(44)}),s_{(44)})\}ds_{(44)}\right]}\ge 0$$

$$T_i(t)\ge T_i^0e^{(-(b_i^{'})^{(9)}t)}>0\quad\text{for t}>0$$

**Definition of** $\left((\widehat{M}_{44})^{(9)}\right)_1,\left((\widehat{M}_{44})^{(9)}\right)_2$ *and* $\left((\widehat{M}_{44})^{(9)}\right)_3$ :

**Remark 43:** if $G_{44}$ is bounded, the same property have also $G_{45}\ and\ G_{46}$ . indeed if

$G_{44} < (\widehat{M}_{44})^{(9)}$ it follows $\frac{dG_{45}}{dt} \leq \left((\widehat{M}_{44})^{(9)}\right)_1 - (a_{45}^{'})^{(9)} G_{45}$ and by integrating

$G_{45} \leq \left((\widehat{M}_{44})^{(9)}\right)_2 = G_{45}^0 + 2(a_{45})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_1/(a_{45}^{'})^{(9)}$

In the same way , one can obtain

$G_{46} \leq \left((\widehat{M}_{44})^{(9)}\right)_3 = G_{46}^0 + 2(a_{46})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_2/(a_{46}^{'})^{(9)}$

If $G_{45}\ or\ G_{46}$ is bounded, the same property follows for $G_{44}$ , $G_{46}$ and $G_{44}$ , $G_{45}$ respectively.

**Remark 44:** If $G_{44}$ *is* bounded, from below, the same property holds for $G_{45}\ and\ G_{46}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{45}$ is bounded from below.

**Remark 45:** If $T_{44}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(9)}\left((G_{47})(t),t\right) = (b_{45}^{'})^{(9)}$ then $T_{45} \to \infty$.

**Definition of** $(m)^{(9)}$ and $\varepsilon_9$ :

Indeed let $t_9$ be so that for $t > t_9$

$$(b_{45})^{(9)} - (b_i^{''})^{(9)}\left((G_{47})(t),t\right) < \varepsilon_9, T_{44}(t) > (m)^{(9)}$$

Then $\frac{dT_{45}}{dt} \geq (a_{45})^{(9)}(m)^{(9)} - \varepsilon_9 T_{45}$ which leads to

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{\varepsilon_9}\right)(1 - e^{-\varepsilon_9 t}) + T_{45}^0 e^{-\varepsilon_9 t}$ If we take t such that $e^{-\varepsilon_9 t} = \frac{1}{2}$ it results

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{2}\right),\quad t = log\frac{2}{\varepsilon_9}$ By taking now $\varepsilon_9$ sufficiently small one sees that $T_{45}$ is unbounded.
The same property holds for $T_{46}$ if $\lim_{t\to\infty}(b_{46}^{''})^{(9)}\left((G_{47})(t),t\right) = (b_{46}^{'})^{(9)}$
We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 92

**Behavior of the solutions of equation** 280

**Theorem** If we denote and define

**Definition of** $(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ :

$(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ four constants satisfying
$-(\sigma_2)^{(1)} \leq -(a_{13}^{'})^{(1)} + (a_{14}^{'})^{(1)} - (a_{13}^{''})^{(1)}(T_{14},t) + (a_{14}^{''})^{(1)}(T_{14},t) \leq -(\sigma_1)^{(1)}$

$-(\tau_2)^{(1)} \leq -(b_{13}^{'})^{(1)} + (b_{14}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G,t) - (b_{14}^{''})^{(1)}(G,t) \leq -(\tau_1)^{(1)}$

**Definition of** $(v_1)^{(1)}, (v_2)^{(1)}, (u_1)^{(1)}, (u_2)^{(1)}, v^{(1)}, u^{(1)}$ : 281

By $(v_1)^{(1)} > 0$ , $(v_2)^{(1)} < 0$ and respectively $(u_1)^{(1)} > 0$ , $(u_2)^{(1)} < 0$ the roots of the equations
$(a_{14})^{(1)}\left(v^{(1)}\right)^2 + (\sigma_1)^{(1)}v^{(1)} - (a_{13})^{(1)} = 0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2 + (\tau_1)^{(1)}u^{(1)} - (b_{13})^{(1)} = 0$

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| **Definition of** $(\bar{v}_1)^{(1)},,(\bar{v}_2)^{(1)},(\bar{u}_1)^{(1)},(\bar{u}_2)^{(1)}$ :<br><br>By $(\bar{v}_1)^{(1)}>0$ , $(\bar{v}_2)^{(1)}<0$ and respectively $(\bar{u}_1)^{(1)}>0$ , $(\bar{u}_2)^{(1)}<0$ the roots of the equations $(a_{14})^{(1)}\left(v^{(1)}\right)^2+(\sigma_2)^{(1)}v^{(1)}-(a_{13})^{(1)}=0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2+(\tau_2)^{(1)}u^{(1)}-(b_{13})^{(1)}=0$ | 282 |
| **Definition of** $(m_1)^{(1)}$ , $(m_2)^{(1)}$ , $(\mu_1)^{(1)},(\mu_2)^{(1)},(\nu_0)^{(1)}$ :-<br><br>If we define $(m_1)^{(1)}$ , $(m_2)^{(1)}$ , $(\mu_1)^{(1)},(\mu_2)^{(1)}$ by<br>$(m_2)^{(1)}=(\nu_0)^{(1)},(m_1)^{(1)}=(\nu_1)^{(1)},\ if\ (\nu_0)^{(1)}<(\nu_1)^{(1)}$<br><br>$(m_2)^{(1)}=(\nu_1)^{(1)},(m_1)^{(1)}=(\bar{\nu}_1)^{(1)}\ ,if\ (\nu_1)^{(1)}<(\nu_0)^{(1)}<(\bar{\nu}_1)^{(1)},$<br><br>and $\boxed{(\nu_0)^{(1)}=\frac{G_{13}^0}{G_{14}^0}}$<br><br>$(\ m_2)^{(1)}=(\nu_1)^{(1)},(m_1)^{(1)}=(\nu_0)^{(1)},\ if\ (\bar{\nu}_1)^{(1)}<(\nu_0)^{(1)}$ | 283 |
| and analogously<br><br>$(\mu_2)^{(1)}=(u_0)^{(1)},(\mu_1)^{(1)}=(u_1)^{(1)},\ if\ (u_0)^{(1)}<(u_1)^{(1)}$<br><br>$(\mu_2)^{(1)}=(u_1)^{(1)},(\mu_1)^{(1)}=(\bar{u}_1)^{(1)}\ ,if\ (u_1)^{(1)}<(u_0)^{(1)}<(\bar{u}_1)^{(1)},$<br><br>and $\boxed{(u_0)^{(1)}=\frac{T_{13}^0}{T_{14}^0}}$<br><br>$(\ \mu_2)^{(1)}=(u_1)^{(1)},(\mu_1)^{(1)}=(u_0)^{(1)},if\ (\bar{u}_1)^{(1)}<(u_0)^{(1)}$ where $(u_1)^{(1)},(\bar{u}_1)^{(1)}$<br><br>are defined | 284 |
| Then the solution of global equations satisfies the inequalities<br><br>$G_{13}^0e^{((S_1)^{(1)}-(p_{13})^{(1)})t}\le G_{13}(t)\le G_{13}^0e^{(S_1)^{(1)}t}$<br><br>where $(p_i)^{(1)}$ is defined by equation<br><br>$\frac{1}{(m_1)^{(1)}}G_{13}^0e^{((S_1)^{(1)}-(p_{13})^{(1)})t}\le G_{14}(t)\le\frac{1}{(m_2)^{(1)}}G_{13}^0e^{(S_1)^{(1)}t}$ | 285 |
| $(\frac{(a_{15})^{(1)}G_{13}^0}{(m_1)^{(1)}((S_1)^{(1)}-(p_{13})^{(1)}-(S_2)^{(1)})}\left[e^{((S_1)^{(1)}-(p_{13})^{(1)})t}-e^{-(S_2)^{(1)}t}\right]+G_{15}^0e^{-(S_2)^{(1)}t}\le G_{15}(t)\le$<br><br>$\frac{(a_{15})^{(1)}G_{13}^0}{(m_2)^{(1)}((S_1)^{(1)}-(a_{15}^{'})^{(1)})}[e^{(S_1)^{(1)}t}-e^{-(a_{15}^{'})^{(1)}t}]+\ G_{15}^0e^{-(a_{15}^{'})^{(1)}t})$ | 286 |
| $\boxed{T_{13}^0e^{(R_1)^{(1)}t}\le T_{13}(t)\le T_{13}^0e^{((R_1)^{(1)}+(r_{13})^{(1)})t}}$ | 287 |
| $\frac{1}{(\mu_1)^{(1)}}T_{13}^0e^{(R_1)^{(1)}t}\le T_{13}(t)\le\frac{1}{(\mu_2)^{(1)}}T_{13}^0e^{((R_1)^{(1)}+(r_{13})^{(1)})t}$ | 288 |

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| $\frac{(b_{15})^{(1)}T_{13}^{0}}{(\mu_1)^{(1)}((R_1)^{(1)}-(b_{15}^{'})^{(1)})}\left[e^{(R_1)^{(1)}t}-e^{-(b_{15}^{'})^{(1)}t}\right]+T_{15}^{0}e^{-(b_{15}^{'})^{(1)}t}\le T_{15}(t)\le$ $\frac{(a_{15})^{(1)}T_{13}^{0}}{(\mu_2)^{(1)}((R_1)^{(1)}+(r_{13})^{(1)}+(R_2)^{(1)})}\left[e^{((R_1)^{(1)}+(r_{13})^{(1)})t}-e^{-(R_2)^{(1)}t}\right]+T_{15}^{0}e^{-(R_2)^{(1)}t}$ | 289 |
| **Definition of** $(S_1)^{(1)},(S_2)^{(1)},(R_1)^{(1)},(R_2)^{(1)}$:- Where $(S_1)^{(1)}=(a_{13})^{(1)}(m_2)^{(1)}-(a_{13}^{'})^{(1)}$ $(S_2)^{(1)}=(a_{15})^{(1)}-(p_{15})^{(1)}$ $(R_1)^{(1)}=(b_{13})^{(1)}(\mu_2)^{(1)}-(b_{13}^{'})^{(1)}$ $(R_2)^{(1)}=(b_{15}^{'})^{(1)}-(r_{15})^{(1)}$ | 290 |
| **Behavior of the solutions of equation** **Theorem 2:** If we denote and define | 291 |
| **Definition of** $(\sigma_1)^{(2)},(\sigma_2)^{(2)},(\tau_1)^{(2)},(\tau_2)^{(2)}$ : $(\sigma_1)^{(2)},(\sigma_2)^{(2)},(\tau_1)^{(2)},(\tau_2)^{(2)}$ four constants satisfying | 292 |
| $-(\sigma_2)^{(2)}\le-(a_{16}^{'})^{(2)}+(a_{17}^{'})^{(2)}-(a_{16}^{''})^{(2)}(\mathrm{T}_{17},t)+(a_{17}^{''})^{(2)}(\mathrm{T}_{17},t)\le-(\sigma_1)^{(2)}$ | 293 |
| $-(\tau_2)^{(2)}\le-(b_{16}^{'})^{(2)}+(b_{17}^{'})^{(2)}-(b_{16}^{''})^{(2)}\big((G_{19}),t\big)-(b_{17}^{''})^{(2)}\big((G_{19}),t\big)\le-(\tau_1)^{(2)}$ | 294 |
| **Definition of** $(v_1)^{(2)},(v_2)^{(2)},(u_1)^{(2)},(u_2)^{(2)}$ : | 295 |
| By $(v_1)^{(2)}>0$ , $(v_2)^{(2)}<0$ and respectively $(u_1)^{(2)}>0$ , $(u_2)^{(2)}<0$ the roots | 296 |
| of the equations $(a_{17})^{(2)}\big(v^{(2)}\big)^2+(\sigma_1)^{(2)}v^{(2)}-(a_{16})^{(2)}=0$ | 297 |
| and $(b_{14})^{(2)}\big(u^{(2)}\big)^2+(\tau_1)^{(2)}u^{(2)}-(b_{16})^{(2)}=0$ and | 298 |
| **Definition of** $(\bar{v}_1)^{(2)},,(\bar{v}_2)^{(2)},(\bar{u}_1)^{(2)},(\bar{u}_2)^{(2)}$ : | 299 |
| By $(\bar{v}_1)^{(2)}>0$ , $(\bar{v}_2)^{(2)}<0$ and respectively $(\bar{u}_1)^{(2)}>0$ , $(\bar{u}_2)^{(2)}<0$ the | 300 |
| roots of the equations $(a_{17})^{(2)}\big(v^{(2)}\big)^2+(\sigma_2)^{(2)}v^{(2)}-(a_{16})^{(2)}=0$ | 301 |
| and $(b_{17})^{(2)}\big(u^{(2)}\big)^2+(\tau_2)^{(2)}u^{(2)}-(b_{16})^{(2)}=0$ | 302 |
| **Definition of** $(m_1)^{(2)},(m_2)^{(2)},(\mu_1)^{(2)},(\mu_2)^{(2)}$ :- | 303 |
| If we define $(m_1)^{(2)},(m_2)^{(2)},(\mu_1)^{(2)},(\mu_2)^{(2)}$ by | 304 |
| $(m_2)^{(2)}=(v_0)^{(2)},(m_1)^{(2)}=(v_1)^{(2)}$, ***if*** $(v_0)^{(2)}<(v_1)^{(2)}$ | 305 |

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| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ , ***if*** $(\nu_1)^{(2)} < (\nu_0)^{(2)} < (\bar{\nu}_1)^{(2)}$, and $\boxed{(\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0}}$ | 306 |
| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\nu_0)^{(2)}$, ***if*** $(\bar{\nu}_1)^{(2)} < (\nu_0)^{(2)}$ | 307 |
| and analogously<br>$(\mu_2)^{(2)} = (u_0)^{(2)}, (\mu_1)^{(2)} = (u_1)^{(2)}$, ***if*** $(u_0)^{(2)} < (u_1)^{(2)}$<br>$(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (\bar{u}_1)^{(2)}$ , ***if*** $(u_1)^{(2)} < (u_0)^{(2)} < (\bar{u}_1)^{(2)}$, and $\boxed{(u_0)^{(2)} = \frac{T_{16}^0}{T_{17}^0}}$ | 308 |
| $(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (u_0)^{(2)}$, ***if*** $(\bar{u}_1)^{(2)} < (u_0)^{(2)}$ | 309 |
| Then the solution of global equations satisfies the inequalities<br>$G_{16}^0 e^{((S_1)^{(2)} - (p_{16})^{(2)})t} \le G_{16}(t) \le G_{16}^0 e^{(S_1)^{(2)} t}$ | 310 |
| $(p_i)^{(2)}$ is defined by equation | |
| $\frac{1}{(m_1)^{(2)}} G_{16}^0 e^{((S_1)^{(2)} - (p_{16})^{(2)})t} \le G_{17}(t) \le \frac{1}{(m_2)^{(2)}} G_{16}^0 e^{(S_1)^{(2)} t}$ | 311 |
| $\left(\frac{(a_{18})^{(2)} G_{16}^0}{(m_1)^{(2)}((S_1)^{(2)} - (p_{16})^{(2)} - (S_2)^{(2)})}\left[e^{((S_1)^{(2)} - (p_{16})^{(2)})t} - e^{-(S_2)^{(2)} t}\right] + G_{18}^0 e^{-(S_2)^{(2)} t} \le G_{18}(t) \le \right.$<br>$\left. \frac{(a_{18})^{(2)} G_{16}^0}{(m_2)^{(2)}((S_1)^{(2)} - (a'_{18})^{(2)})}\left[e^{(S_1)^{(2)} t} - e^{-(a'_{18})^{(2)} t}\right] + G_{18}^0 e^{-(a'_{18})^{(2)} t}\right)$ | 312 |
| $\boxed{T_{16}^0 e^{(R_1)^{(2)} t} \le T_{16}(t) \le T_{16}^0 e^{((R_1)^{(2)} + (r_{16})^{(2)})t}}$ | 313 |
| $\frac{1}{(\mu_1)^{(2)}} T_{16}^0 e^{(R_1)^{(2)} t} \le T_{16}(t) \le \frac{1}{(\mu_2)^{(2)}} T_{16}^0 e^{((R_1)^{(2)} + (r_{16})^{(2)})t}$ | 314 |
| $\frac{(b_{18})^{(2)} T_{16}^0}{(\mu_1)^{(2)}((R_1)^{(2)} - (b'_{18})^{(2)})}\left[e^{(R_1)^{(2)} t} - e^{-(b'_{18})^{(2)} t}\right] + T_{18}^0 e^{-(b'_{18})^{(2)} t} \le T_{18}(t) \le$<br>$\frac{(a_{18})^{(2)} T_{16}^0}{(\mu_2)^{(2)}((R_1)^{(2)} + (r_{16})^{(2)} + (R_2)^{(2)})}\left[e^{((R_1)^{(2)} + (r_{16})^{(2)})t} - e^{-(R_2)^{(2)} t}\right] + T_{18}^0 e^{-(R_2)^{(2)} t}$ | 315 |
| **Definition of** $(S_1)^{(2)}, (S_2)^{(2)}, (R_1)^{(2)}, (R_2)^{(2)}$:- | 316 |
| Where $(S_1)^{(2)} = (a_{16})^{(2)}(m_2)^{(2)} - (a'_{16})^{(2)}$<br>$(S_2)^{(2)} = (a_{18})^{(2)} - (p_{18})^{(2)}$ | 317 |
| $(R_1)^{(2)} = (b_{16})^{(2)}(\mu_2)^{(1)} - (b'_{16})^{(2)}$ | 318 |

$(R_2)^{(2)} = (b'_{18})^{(2)} - (r_{18})^{(2)}$

**Behavior of the solutions** 319

**Theorem 3:** If we denote and define

**Definition of** $(\sigma_1)^{(3)}, (\sigma_2)^{(3)}, (\tau_1)^{(3)}, (\tau_2)^{(3)}$ :

$(\sigma_1)^{(3)}, (\sigma_2)^{(3)}, (\tau_1)^{(3)}, (\tau_2)^{(3)}$ four constants satisfying

$$-(\sigma_2)^{(3)} \le -(a'_{20})^{(3)} + (a'_{21})^{(3)} - (a''_{20})^{(3)}(T_{21}, t) + (a''_{21})^{(3)}(T_{21}, t) \le -(\sigma_1)^{(3)}$$

$$-(\tau_2)^{(3)} \le -(b'_{20})^{(3)} + (b'_{21})^{(3)} - (b''_{20})^{(3)}(G_{23}, t) - (b''_{21})^{(3)}\big((G_{23}), t\big) \le -(\tau_1)^{(3)}$$

**Definition of** $(v_1)^{(3)}, (v_2)^{(3)}, (u_1)^{(3)}, (u_2)^{(3)}$ : 320

By $(v_1)^{(3)} > 0$ , $(v_2)^{(3)} < 0$ and respectively $(u_1)^{(3)} > 0$ , $(u_2)^{(3)} < 0$ the roots of the equations $(a_{21})^{(3)}\big(v^{(3)}\big)^2 + (\sigma_1)^{(3)}v^{(3)} - (a_{20})^{(3)} = 0$

and $(b_{21})^{(3)}\big(u^{(3)}\big)^2 + (\tau_1)^{(3)}u^{(3)} - (b_{20})^{(3)} = 0$ and

By $(\bar{v}_1)^{(3)} > 0$ , $(\bar{v}_2)^{(3)} < 0$ and respectively $(\bar{u}_1)^{(3)} > 0$ , $(\bar{u}_2)^{(3)} < 0$ the

roots of the equations $(a_{21})^{(3)}\big(v^{(3)}\big)^2 + (\sigma_2)^{(3)}v^{(3)} - (a_{20})^{(3)} = 0$

and $(b_{21})^{(3)}\big(u^{(3)}\big)^2 + (\tau_2)^{(3)}u^{(3)} - (b_{20})^{(3)} = 0$

**Definition of** $(m_1)^{(3)}, (m_2)^{(3)}, (\mu_1)^{(3)}, (\mu_2)^{(3)}$ :- 321

If we define $(m_1)^{(3)}, (m_2)^{(3)}, (\mu_1)^{(3)}, (\mu_2)^{(3)}$ by

$(m_2)^{(3)} = (v_0)^{(3)}, (m_1)^{(3)} = (v_1)^{(3)}$, ***if*** $(v_0)^{(3)} < (v_1)^{(3)}$

$(m_2)^{(3)} = (v_1)^{(3)}, (m_1)^{(3)} = (\bar{v}_1)^{(3)}$ , ***if*** $(v_1)^{(3)} < (v_0)^{(3)} < (\bar{v}_1)^{(3)}$,

and $\boxed{(v_0)^{(3)} = \frac{G^0_{20}}{G^0_{21}}}$

$(m_2)^{(3)} = (v_1)^{(3)}, (m_1)^{(3)} = (v_0)^{(3)}$, ***if*** $(\bar{v}_1)^{(3)} < (v_0)^{(3)}$

and analogously 322

$(\mu_2)^{(3)} = (u_0)^{(3)}, (\mu_1)^{(3)} = (u_1)^{(3)}$, ***if*** $(u_0)^{(3)} < (u_1)^{(3)}$

$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (\bar{u}_1)^{(3)}$ , ***if*** $(u_1)^{(3)} < (u_0)^{(3)} < (\bar{u}_1)^{(3)}$, and $\boxed{(u_0)^{(3)} = \frac{T^0_{20}}{T^0_{21}}}$

$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (u_0)^{(3)}$, ***if*** $(\bar{u}_1)^{(3)} < (u_0)^{(3)}$

Then the solution of global equations satisfies the inequalities

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| $G_{20}^0 e^{((S_1)^{(3)}-(p_{20})^{(3)})t} \le G_{20}(t) \le G_{20}^0 e^{(S_1)^{(3)}t}$<br>$(p_i)^{(3)}$ is defined by equation | |
| $\frac{1}{(m_1)^{(3)}} G_{20}^0 e^{((S_1)^{(3)}-(p_{20})^{(3)})t} \le G_{21}(t) \le \frac{1}{(m_2)^{(3)}} G_{20}^0 e^{(S_1)^{(3)}t}$ | 323 |
| $\left(\frac{(a_{22})^{(3)}G_{20}^0}{(m_1)^{(3)}((S_1)^{(3)}-(p_{20})^{(3)}-(S_2)^{(3)})}\left[e^{((S_1)^{(3)}-(p_{20})^{(3)})t} - e^{-(S_2)^{(3)}t}\right] + G_{22}^0 e^{-(S_2)^{(3)}t} \le G_{22}(t) \le \right.$<br>$\left.\frac{(a_{22})^{(3)}G_{20}^0}{(m_2)^{(3)}((S_1)^{(3)}-(a'_{22})^{(3)})}[e^{(S_1)^{(3)}t} - e^{-(a'_{22})^{(3)}t}] + G_{22}^0 e^{-(a'_{22})^{(3)}t}\right)$ | 324 |
| $\boxed{T_{20}^0 e^{(R_1)^{(3)}t} \le T_{20}(t) \le T_{20}^0 e^{((R_1)^{(3)}+(r_{20})^{(3)})t}}$ | 325 |
| $\frac{1}{(\mu_1)^{(3)}} T_{20}^0 e^{(R_1)^{(3)}t} \le T_{20}(t) \le \frac{1}{(\mu_2)^{(3)}} T_{20}^0 e^{((R_1)^{(3)}+(r_{20})^{(3)})t}$ | 326 |
| $\frac{(b_{22})^{(3)}T_{20}^0}{(\mu_1)^{(3)}((R_1)^{(3)}-(b'_{22})^{(3)})}\left[e^{(R_1)^{(3)}t} - e^{-(b'_{22})^{(3)}t}\right] + T_{22}^0 e^{-(b'_{22})^{(3)}t} \le T_{22}(t) \le$<br>$\frac{(a_{22})^{(3)}T_{20}^0}{(\mu_2)^{(3)}((R_1)^{(3)}+(r_{20})^{(3)}+(R_2)^{(3)})}\left[e^{((R_1)^{(3)}+(r_{20})^{(3)})t} - e^{-(R_2)^{(3)}t}\right] + T_{22}^0 e^{-(R_2)^{(3)}t}$ | 327 |
| **Definition of** $(S_1)^{(3)}, (S_2)^{(3)}, (R_1)^{(3)}, (R_2)^{(3)}$:-<br>Where $(S_1)^{(3)} = (a_{20})^{(3)}(m_2)^{(3)} - (a'_{20})^{(3)}$<br>$(S_2)^{(3)} = (a_{22})^{(3)} - (p_{22})^{(3)}$<br>$(R_1)^{(3)} = (b_{20})^{(3)}(\mu_2)^{(3)} - (b'_{20})^{(3)}$<br>$(R_2)^{(3)} = (b'_{22})^{(3)} - (r_{22})^{(3)}$ | 328 |
| **Behavior of the solutions of equation**<br>**Theorem:** If we denote and define<br>**Definition of** $(\sigma_1)^{(4)}, (\sigma_2)^{(4)}, (\tau_1)^{(4)}, (\tau_2)^{(4)}$ :<br>$(\sigma_1)^{(4)}, (\sigma_2)^{(4)}, (\tau_1)^{(4)}, (\tau_2)^{(4)}$ four constants satisfying<br>$-(\sigma_2)^{(4)} \le -(a'_{24})^{(4)} + (a'_{25})^{(4)} - (a''_{24})^{(4)}(T_{25}, t) + (a''_{25})^{(4)}(T_{25}, t) \le -(\sigma_1)^{(4)}$<br>$-(\tau_2)^{(4)} \le -(b'_{24})^{(4)} + (b'_{25})^{(4)} - (b''_{24})^{(4)}((G_{27}), t) - (b''_{25})^{(4)}((G_{27}), t) \le -(\tau_1)^{(4)}$ | |
| **Definition of** $(v_1)^{(4)}, (v_2)^{(4)}, (u_1)^{(4)}, (u_2)^{(4)}, v^{(4)}, u^{(4)}$ :<br>By $(v_1)^{(4)} > 0$, $(v_2)^{(4)} < 0$ and respectively $(u_1)^{(4)} > 0$, $(u_2)^{(4)} < 0$ the roots of the equations $(a_{25})^{(4)}(v^{(4)})^2 + (\sigma_1)^{(4)}v^{(4)} - (a_{24})^{(4)} = 0$<br>and $(b_{25})^{(4)}(u^{(4)})^2 + (\tau_1)^{(4)}u^{(4)} - (b_{24})^{(4)} = 0$ and | 329 |

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| **Definition of** $(\bar{v}_1)^{(4)},,(\bar{v}_2)^{(4)},(\bar{u}_1)^{(4)},(\bar{u}_2)^{(4)}$ :<br><br>By $(\bar{v}_1)^{(4)}>0$ , $(\bar{v}_2)^{(4)}<0$ and respectively $(\bar{u}_1)^{(4)}>0$ , $(\bar{u}_2)^{(4)}<0$ the roots of the equations $(a_{25})^{(4)}\left(v^{(4)}\right)^2+(\sigma_2)^{(4)}v^{(4)}-(a_{24})^{(4)}=0$ and $(b_{25})^{(4)}\left(u^{(4)}\right)^2+(\tau_2)^{(4)}u^{(4)}-(b_{24})^{(4)}=0$<br>**Definition of** $(m_1)^{(4)}$ , $(m_2)^{(4)}$ , $(\mu_1)^{(4)},(\mu_2)^{(4)},(v_0)^{(4)}$ :-<br><br>If we define $(m_1)^{(4)}$ , $(m_2)^{(4)}$ , $(\mu_1)^{(4)},(\mu_2)^{(4)}$ by<br><br>$(m_2)^{(4)}=(v_0)^{(4)},(m_1)^{(4)}=(v_1)^{(4)}$, ***if*** $(v_0)^{(4)}<(v_1)^{(4)}$<br><br>$(m_2)^{(4)}=(v_1)^{(4)},(m_1)^{(4)}=(\bar{v}_1)^{(4)}$ , ***if*** $(v_4)^{(4)}<(v_0)^{(4)}<(\bar{v}_1)^{(4)}$,<br>and $\boxed{(v_0)^{(4)}=\frac{G_{24}^0}{G_{25}^0}}$<br><br>$(m_2)^{(4)}=(v_4)^{(4)},(m_1)^{(4)}=(v_0)^{(4)}$, ***if*** $(\bar{v}_4)^{(4)}<(v_0)^{(4)}$ | 330 |
| and analogously<br><br>$(\mu_2)^{(4)}=(u_0)^{(4)},(\mu_1)^{(4)}=(u_1)^{(4)}$, ***if*** $(u_0)^{(4)}<(u_1)^{(4)}$<br><br>$(\mu_2)^{(4)}=(u_1)^{(4)},(\mu_1)^{(4)}=(\bar{u}_1)^{(4)}$ , ***if*** $(u_1)^{(4)}<(u_0)^{(4)}<(\bar{u}_1)^{(4)}$,<br>and $\boxed{(u_0)^{(4)}=\frac{T_{24}^0}{T_{25}^0}}$<br><br>$(\mu_2)^{(4)}=(u_1)^{(4)},(\mu_1)^{(4)}=(u_0)^{(4)}$, ***if*** $(\bar{u}_1)^{(4)}<(u_0)^{(4)}$ where $(u_1)^{(4)},(\bar{u}_1)^{(4)}$ | 331 |
| Then the solution of global equations satisfies the inequalities<br><br>$G_{24}^0e^{((S_1)^{(4)}-(p_{24})^{(4)})t}\le G_{24}(t)\le G_{24}^0e^{(S_1)^{(4)}t}$<br><br>where $(p_i)^{(4)}$ is defined by equation | 332 |
| $\frac{1}{(m_1)^{(4)}}G_{24}^0e^{((S_1)^{(4)}-(p_{24})^{(4)})t}\le G_{25}(t)\le\frac{1}{(m_2)^{(4)}}G_{24}^0e^{(S_1)^{(4)}t}$ | 333 |
| $\left(\frac{(a_{26})^{(4)}G_{24}^0}{(m_1)^{(4)}((S_1)^{(4)}-(p_{24})^{(4)}-(S_2)^{(4)})}\left[e^{((S_1)^{(4)}-(p_{24})^{(4)})t}-e^{-(S_2)^{(4)}t}\right]+G_{26}^0e^{-(S_2)^{(4)}t}\le G_{26}(t)\le\right.$<br>*(a26)4G240(m2)4(S1)4−(a26′)4e(S1)4t−e−(a26′)4t+ G260e−(a26′)4t* | 334 |
| $\boxed{T_{24}^0e^{(R_1)^{(4)}t}\le T_{24}(t)\le T_{24}^0e^{((R_1)^{(4)}+(r_{24})^{(4)})t}}$ | |
| $\frac{1}{(\mu_1)^{(4)}}T_{24}^0e^{(R_1)^{(4)}t}\le T_{24}(t)\le\frac{1}{(\mu_2)^{(4)}}T_{24}^0e^{((R_1)^{(4)}+(r_{24})^{(4)})t}$ | 335 |
| $\frac{(b_{26})^{(4)}T_{24}^0}{(\mu_1)^{(4)}((R_1)^{(4)}-(b_{26}^{'})^{(4)})}\left[e^{(R_1)^{(4)}t}-e^{-(b_{26}^{'})^{(4)}t}\right]+T_{26}^0e^{-(b_{26}^{'})^{(4)}t}\le T_{26}(t)\le$<br><br>$\frac{(a_{26})^{(4)}T_{24}^0}{(\mu_2)^{(4)}((R_1)^{(4)}+(r_{24})^{(4)}+(R_2)^{(4)})}\left[e^{((R_1)^{(4)}+(r_{24})^{(4)})t}-e^{-(R_2)^{(4)}t}\right]+T_{26}^0e^{-(R_2)^{(4)}t}$ | 336 |

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| **Definition of** $(S_1)^{(4)}, (S_2)^{(4)}, (R_1)^{(4)}, (R_2)^{(4)}$:-<br>Where $(S_1)^{(4)} = (a_{24})^{(4)}(m_2)^{(4)} - (a'_{24})^{(4)}$<br>$(S_2)^{(4)} = (a_{26})^{(4)} - (p_{26})^{(4)}$<br>$(R_1)^{(4)} = (b_{24})^{(4)}(\mu_2)^{(4)} - (b'_{24})^{(4)}$<br>$(R_2)^{(4)} = (b'_{26})^{(4)} - (r_{26})^{(4)}$ | 337 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ :<br>$(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ four constants satisfying<br>$-(\sigma_2)^{(5)} \le -(a'_{28})^{(5)} + (a'_{29})^{(5)} - (a''_{28})^{(5)}(T_{29}, t) + (a''_{29})^{(5)}(T_{29}, t) \le -(\sigma_1)^{(5)}$<br>$-(\tau_2)^{(5)} \le -(b'_{28})^{(5)} + (b'_{29})^{(5)} - (b''_{28})^{(5)}((G_{31}), t) - (b''_{29})^{(5)}((G_{31}), t) \le -(\tau_1)^{(5)}$ | 338 |
| **Definition of** $(v_1)^{(5)}, (v_2)^{(5)}, (u_1)^{(5)}, (u_2)^{(5)}, v^{(5)}, u^{(5)}$ :<br>By $(v_1)^{(5)} > 0$, $(v_2)^{(5)} < 0$ and respectively $(u_1)^{(5)} > 0$, $(u_2)^{(5)} < 0$ the roots of the equations $(a_{29})^{(5)}\left(v^{(5)}\right)^2 + (\sigma_1)^{(5)}v^{(5)} - (a_{28})^{(5)} = 0$<br>and $(b_{29})^{(5)}\left(u^{(5)}\right)^2 + (\tau_1)^{(5)}u^{(5)} - (b_{28})^{(5)} = 0$ and | 339 |
| **Definition of** $(\bar{v}_1)^{(5)},, (\bar{v}_2)^{(5)}, (\bar{u}_1)^{(5)}, (\bar{u}_2)^{(5)}$ :<br>By $(\bar{v}_1)^{(5)} > 0$, $(\bar{v}_2)^{(5)} < 0$ and respectively $(\bar{u}_1)^{(5)} > 0$, $(\bar{u}_2)^{(5)} < 0$ the roots of the equations $(a_{29})^{(5)}\left(v^{(5)}\right)^2 + (\sigma_2)^{(5)}v^{(5)} - (a_{28})^{(5)} = 0$<br>and $(b_{29})^{(5)}\left(u^{(5)}\right)^2 + (\tau_2)^{(5)}u^{(5)} - (b_{28})^{(5)} = 0$<br>**Definition of** $(m_1)^{(5)}, (m_2)^{(5)}, (\mu_1)^{(5)}, (\mu_2)^{(5)}, (v_0)^{(5)}$ :-<br>If we define $(m_1)^{(5)}, (m_2)^{(5)}, (\mu_1)^{(5)}, (\mu_2)^{(5)}$ by<br>$(m_2)^{(5)} = (v_0)^{(5)}, (m_1)^{(5)} = (v_1)^{(5)}$, ***if*** $(v_0)^{(5)} < (v_1)^{(5)}$<br>$(m_2)^{(5)} = (v_1)^{(5)}, (m_1)^{(5)} = (\bar{v}_1)^{(5)}$, ***if*** $(v_1)^{(5)} < (v_0)^{(5)} < (\bar{v}_1)^{(5)}$,<br>and $\boxed{(v_0)^{(5)} = \frac{G^0_{28}}{G^0_{29}}}$<br>$(m_2)^{(5)} = (v_1)^{(5)}, (m_1)^{(5)} = (v_0)^{(5)}$, ***if*** $(\bar{v}_1)^{(5)} < (v_0)^{(5)}$ | 340 |
| and analogously<br>$(\mu_2)^{(5)} = (u_0)^{(5)}, (\mu_1)^{(5)} = (u_1)^{(5)}$, ***if*** $(u_0)^{(5)} < (u_1)^{(5)}$ | 341 |

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| $(\mu_2)^{(5)} = (u_1)^{(5)}, (\mu_1)^{(5)} = (\bar{u}_1)^{(5)}$ , $\boldsymbol{if}$ $(u_1)^{(5)} < (u_0)^{(5)} < (\bar{u}_1)^{(5)}$, and $\boxed{(u_0)^{(5)} = \frac{T_{28}^0}{T_{29}^0}}$ $(\mu_2)^{(5)} = (u_1)^{(5)}, (\mu_1)^{(5)} = (u_0)^{(5)}, \boldsymbol{if}$ $(\bar{u}_1)^{(5)} < (u_0)^{(5)}$ where $(u_1)^{(5)}, (\bar{u}_1)^{(5)}$ | |
| Then the solution of global equations satisfies the inequalities $G_{28}^0 e^{((S_1)^{(5)} - (p_{28})^{(5)})t} \le G_{28}(t) \le G_{28}^0 e^{(S_1)^{(5)} t}$ where $(p_i)^{(5)}$ is defined by equation | 342 |
| $\frac{1}{(m_5)^{(5)}} G_{28}^0 e^{((S_1)^{(5)} - (p_{28})^{(5)})t} \le G_{29}(t) \le \frac{1}{(m_2)^{(5)}} G_{28}^0 e^{(S_1)^{(5)} t}$ | 343 |
| $\left(\frac{(a_{30})^{(5)} G_{28}^0}{(m_1)^{(5)}((S_1)^{(5)} - (p_{28})^{(5)} - (S_2)^{(5)})}\left[e^{((S_1)^{(5)} - (p_{28})^{(5)})t} - e^{-(S_2)^{(5)} t}\right] + G_{30}^0 e^{-(S_2)^{(5)} t} \le G_{30}(t) \le$ *(a30)5G280(m2)5(S1)5−(a30′)5e(S1)5t−e−(a30′)5t+ G300e−(a30′)5t* | 344 |
| $\boxed{T_{28}^0 e^{(R_1)^{(5)} t} \le T_{28}(t) \le T_{28}^0 e^{((R_1)^{(5)} + (r_{28})^{(5)})t}}$ | 345 |
| $\frac{1}{(\mu_1)^{(5)}} T_{28}^0 e^{(R_1)^{(5)} t} \le T_{28}(t) \le \frac{1}{(\mu_2)^{(5)}} T_{28}^0 e^{((R_1)^{(5)} + (r_{28})^{(5)})t}$ | 346 |
| $\frac{(b_{30})^{(5)} T_{28}^0}{(\mu_1)^{(5)}((R_1)^{(5)} - (b_{30}^{'})^{(5)})}\left[e^{(R_1)^{(5)} t} - e^{-(b_{30}^{'})^{(5)} t}\right] + T_{30}^0 e^{-(b_{30}^{'})^{(5)} t} \le T_{30}(t) \le$ $\frac{(a_{30})^{(5)} T_{28}^0}{(\mu_2)^{(5)}((R_1)^{(5)} + (r_{28})^{(5)} + (R_2)^{(5)})}\left[e^{((R_1)^{(5)} + (r_{28})^{(5)})t} - e^{-(R_2)^{(5)} t}\right] + T_{30}^0 e^{-(R_2)^{(5)} t}$ | 347 |
| **Definition of** $(S_1)^{(5)}, (S_2)^{(5)}, (R_1)^{(5)}, (R_2)^{(5)}$:- Where $(S_1)^{(5)} = (a_{28})^{(5)}(m_2)^{(5)} - (a_{28}^{'})^{(5)}$ $(S_2)^{(5)} = (a_{30})^{(5)} - (p_{30})^{(5)}$ $(R_1)^{(5)} = (b_{28})^{(5)}(\mu_2)^{(5)} - (b_{28}^{'})^{(5)}$ $(R_2)^{(5)} = (b_{30}^{'})^{(5)} - (r_{30})^{(5)}$ | 348 |
| **Behavior of the solutions of equation** **Theorem 2:** If we denote and define **Definition of** $(\sigma_1)^{(6)}, (\sigma_2)^{(6)}, (\tau_1)^{(6)}, (\tau_2)^{(6)}$ : $(\sigma_1)^{(6)}, (\sigma_2)^{(6)}, (\tau_1)^{(6)}, (\tau_2)^{(6)}$ four constants satisfying $-(\sigma_2)^{(6)} \le -(a_{32}^{'})^{(6)} + (a_{33}^{'})^{(6)} - (a_{32}^{''})^{(6)}(T_{33}, t) + (a_{33}^{''})^{(6)}(T_{33}, t) \le -(\sigma_1)^{(6)}$ | 349 |

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| $-(\tau_2)^{(6)} \le -(b'_{32})^{(6)} + (b'_{33})^{(6)} - (b''_{32})^{(6)}\big((G_{35}),t\big) - (b''_{33})^{(6)}\big((G_{35}),t\big) \le -(\tau_1)^{(6)}$ | |
| **Definition of** $(\nu_1)^{(6)}, (\nu_2)^{(6)}, (u_1)^{(6)}, (u_2)^{(6)}, \nu^{(6)}, u^{(6)}$ :<br><br>By $(\nu_1)^{(6)} > 0$ , $(\nu_2)^{(6)} < 0$ and respectively $(u_1)^{(6)} > 0$ , $(u_2)^{(6)} < 0$ the roots of the equations $(a_{33})^{(6)}\big(\nu^{(6)}\big)^2 + (\sigma_1)^{(6)}\nu^{(6)} - (a_{32})^{(6)} = 0$<br>and $(b_{33})^{(6)}\big(u^{(6)}\big)^2 + (\tau_1)^{(6)}u^{(6)} - (b_{32})^{(6)} = 0$ and | 350 |
| **Definition of** $(\bar{\nu}_1)^{(6)},, (\bar{\nu}_2)^{(6)}, (\bar{u}_1)^{(6)}, (\bar{u}_2)^{(6)}$ :<br><br>By $(\bar{\nu}_1)^{(6)} > 0$ , $(\bar{\nu}_2)^{(6)} < 0$ and respectively $(\bar{u}_1)^{(6)} > 0$ , $(\bar{u}_2)^{(6)} < 0$ the roots of the equations $(a_{33})^{(6)}\big(\nu^{(6)}\big)^2 + (\sigma_2)^{(6)}\nu^{(6)} - (a_{32})^{(6)} = 0$<br>and $(b_{33})^{(6)}\big(u^{(6)}\big)^2 + (\tau_2)^{(6)}u^{(6)} - (b_{32})^{(6)} = 0$<br>**Definition of** $(m_1)^{(6)}$ , $(m_2)^{(6)}$ , $(\mu_1)^{(6)}, (\mu_2)^{(6)}, (\nu_0)^{(6)}$ :-<br><br>If we define $(m_1)^{(6)}$ , $(m_2)^{(6)}$ , $(\mu_1)^{(6)}, (\mu_2)^{(6)}$ by<br><br>$(m_2)^{(6)} = (\nu_0)^{(6)}, (m_1)^{(6)} = (\nu_1)^{(6)}$, ***if*** $(\nu_0)^{(6)} < (\nu_1)^{(6)}$<br><br>$(m_2)^{(6)} = (\nu_1)^{(6)}, (m_1)^{(6)} = (\bar{\nu}_6)^{(6)}$ , ***if*** $(\nu_1)^{(6)} < (\nu_0)^{(6)} < (\bar{\nu}_1)^{(6)}$,<br>and $\boxed{(\nu_0)^{(6)} = \frac{G^0_{32}}{G^0_{33}}}$<br><br>$(m_2)^{(6)} = (\nu_1)^{(6)}, (m_1)^{(6)} = (\nu_0)^{(6)}$, ***if*** $(\bar{\nu}_1)^{(6)} < (\nu_0)^{(6)}$ | 351 |
| and analogously<br><br>$(\mu_2)^{(6)} = (u_0)^{(6)}, (\mu_1)^{(6)} = (u_1)^{(6)}$, ***if*** $(u_0)^{(6)} < (u_1)^{(6)}$<br><br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (\bar{u}_1)^{(6)}$ , ***if*** $(u_1)^{(6)} < (u_0)^{(6)} < (\bar{u}_1)^{(6)}$,<br>and $\boxed{(u_0)^{(6)} = \frac{T^0_{32}}{T^0_{33}}}$<br><br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (u_0)^{(6)}$, ***if*** $(\bar{u}_1)^{(6)} < (u_0)^{(6)}$ where $(u_1)^{(6)}, (\bar{u}_1)^{(6)}$ | 352 |
| Then the solution of global equations satisfies the inequalities<br><br>$G^0_{32}e^{((S_1)^{(6)} - (p_{32})^{(6)})t} \le G_{32}(t) \le G^0_{32}e^{(S_1)^{(6)}t}$<br><br>where $(p_i)^{(6)}$ is defined by equation | 353 |
| $\frac{1}{(m_1)^{(6)}} G^0_{32}e^{((S_1)^{(6)} - (p_{32})^{(6)})t} \le G_{33}(t) \le \frac{1}{(m_2)^{(6)}} G^0_{32}e^{(S_1)^{(6)}t}$ | 354 |
| $\left(\frac{(a_{34})^{(6)}G^0_{32}}{(m_1)^{(6)}((S_1)^{(6)} - (p_{32})^{(6)} - (S_2)^{(6)})}\left[e^{((S_1)^{(6)} - (p_{32})^{(6)})t} - e^{-(S_2)^{(6)}t}\right] + G^0_{34}e^{-(S_2)^{(6)}t} \le G_{34}(t) \le$<br>*(a34)6G320(m2)6(S1)6−(a34′)6e(S1)6t−e−(a34′)6t+ G340e−(a34′)6t* | 355 |

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| $\boxed{T_{32}^0 e^{(R_1)^{(6)}t} \le T_{32}(t) \le T_{32}^0 e^{((R_1)^{(6)}+(r_{32})^{(6)})t}}$ | 356 |
| $\frac{1}{(\mu_1)^{(6)}} T_{32}^0 e^{(R_1)^{(6)}t} \le T_{32}(t) \le \frac{1}{(\mu_2)^{(6)}} T_{32}^0 e^{((R_1)^{(6)}+(r_{32})^{(6)})t}$ | 357 |
| $\frac{(b_{34})^{(6)}T_{32}^0}{(\mu_1)^{(6)}((R_1)^{(6)}-(b_{34}')^{(6)})}\left[e^{(R_1)^{(6)}t} - e^{-(b_{34}')^{(6)}t}\right] + T_{34}^0 e^{-(b_{34}')^{(6)}t} \le T_{34}(t) \le$ $\frac{(a_{34})^{(6)}T_{32}^0}{(\mu_2)^{(6)}((R_1)^{(6)}+(r_{32})^{(6)}+(R_2)^{(6)})}\left[e^{((R_1)^{(6)}+(r_{32})^{(6)})t} - e^{-(R_2)^{(6)}t}\right] + T_{34}^0 e^{-(R_2)^{(6)}t}$ | 358 |
| **Definition of** $(S_1)^{(6)}, (S_2)^{(6)}, (R_1)^{(6)}, (R_2)^{(6)}$:- Where $(S_1)^{(6)} = (a_{32})^{(6)}(m_2)^{(6)} - (a_{32}')^{(6)}$ $(S_2)^{(6)} = (a_{34})^{(6)} - (p_{34})^{(6)}$ $(R_1)^{(6)} = (b_{32})^{(6)}(\mu_2)^{(6)} - (b_{32}')^{(6)}$ $(R_2)^{(6)} = (b_{34}')^{(6)} - (r_{34})^{(6)}$ | 359 |
| **Behavior of the solutions of equation**<br><br>**Theorem 2:** If we denote and define<br><br>**Definition of** $(\sigma_1)^{(7)}, (\sigma_2)^{(7)}, (\tau_1)^{(7)}, (\tau_2)^{(7)}$ :<br><br>$(\sigma_1)^{(7)}, (\sigma_2)^{(7)}, (\tau_1)^{(7)}, (\tau_2)^{(7)}$ four constants satisfying<br><br>$-(\sigma_2)^{(7)} \le -(a_{36}')^{(7)} + (a_{37}')^{(7)} - (a_{36}'')^{(7)}(T_{37}, t) + (a_{37}'')^{(7)}(T_{37}, t) \le -(\sigma_1)^{(7)}$<br><br>$-(\tau_2)^{(7)} \le -(b_{36}')^{(7)} + (b_{37}')^{(7)} - (b_{36}'')^{(7)}((G_{39}), t) - (b_{37}'')^{(7)}((G_{39}), t) \le -(\tau_1)^{(7)}$ | |
| **Definition of** $(v_1)^{(7)}, (v_2)^{(7)}, (u_1)^{(7)}, (u_2)^{(7)}, v^{(7)}, u^{(7)}$ :<br><br>By $(v_1)^{(7)} > 0, (v_2)^{(7)} < 0$ and respectively $(u_1)^{(7)} > 0, (u_2)^{(7)} < 0$ the roots of the equations $(a_{37})^{(7)}(v^{(7)})^2 + (\sigma_1)^{(7)}v^{(7)} - (a_{36})^{(7)} = 0$ and $(b_{37})^{(7)}(u^{(7)})^2 + (\tau_1)^{(7)}u^{(7)} - (b_{36})^{(7)} = 0$ and | 361 |
| **Definition of** $(\bar{v}_1)^{(7)}, , (\bar{v}_2)^{(7)}, (\bar{u}_1)^{(7)}, (\bar{u}_2)^{(7)}$ :<br><br>By $(\bar{v}_1)^{(7)} > 0, (\bar{v}_2)^{(7)} < 0$ and respectively $(\bar{u}_1)^{(7)} > 0, (\bar{u}_2)^{(7)} < 0$ the roots of the equations $(a_{37})^{(7)}(v^{(7)})^2 + (\sigma_2)^{(7)}v^{(7)} - (a_{36})^{(7)} = 0$ and $(b_{37})^{(7)}(u^{(7)})^2 + (\tau_2)^{(7)}u^{(7)} - (b_{36})^{(7)} = 0$<br>**Definition of** $(m_1)^{(7)}, (m_2)^{(7)}, (\mu_1)^{(7)}, (\mu_2)^{(7)}, (v_0)^{(7)}$ :- | 362 |

If we define $(m_1)^{(7)}, (m_2)^{(7)}, (\mu_1)^{(7)}, (\mu_2)^{(7)}$ by

$$(m_2)^{(7)} = (\nu_0)^{(7)}, (m_1)^{(7)} = (\nu_1)^{(7)}, \ \boldsymbol{if}\ (\nu_0)^{(7)} < (\nu_1)^{(7)}$$

$$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\bar{\nu}_1)^{(7)}, \boldsymbol{if}\ (\nu_1)^{(7)} < (\nu_0)^{(7)} < (\bar{\nu}_1)^{(7)},$$

and $\boxed{(\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0}}$

$$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\nu_0)^{(7)}, \ \boldsymbol{if}\ (\bar{\nu}_1)^{(7)} < (\nu_0)^{(7)}$$

and analogously

$$(\mu_2)^{(7)} = (u_0)^{(7)}, (\mu_1)^{(7)} = (u_1)^{(7)}, \ \boldsymbol{if}\ (u_0)^{(7)} < (u_1)^{(7)}$$

$$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (\bar{u}_1)^{(7)}, \boldsymbol{if}\ (u_1)^{(7)} < (u_0)^{(7)} < (\bar{u}_1)^{(7)},$$

and $\boxed{(u_0)^{(7)} = \frac{T_{36}^0}{T_{37}^0}}$

$$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (u_0)^{(7)}, \boldsymbol{if}\ (\bar{u}_1)^{(7)} < (u_0)^{(7)}$$ where $(u_1)^{(7)}, (\bar{u}_1)^{(7)}$

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Then the solution of global equations satisfies the inequalities

$$G_{36}^0 e^{((S_1)^{(7)} - (p_{36})^{(7)})t} \le G_{36}(t) \le G_{36}^0 e^{(S_1)^{(7)} t}$$

where $(p_i)^{(7)}$ is defined by equation

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$$\frac{1}{(m_7)^{(7)}} G_{36}^0 e^{((S_1)^{(7)} - (p_{36})^{(7)})t} \le G_{37}(t) \le \frac{1}{(m_2)^{(7)}} G_{36}^0 e^{(S_1)^{(7)} t}$$

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$$\left(\frac{(a_{38})^{(7)} G_{36}^0}{(m_1)^{(7)}((S_1)^{(7)} - (p_{36})^{(7)} - (S_2)^{(7)})}\left[e^{((S_1)^{(7)} - (p_{36})^{(7)})t} - e^{-(S_2)^{(7)} t}\right] + G_{38}^0 e^{-(S_2)^{(7)} t} \le G_{38}(t) \le \right.$$

$$\left. \frac{(a_{38})^{(7)} G_{36}^0}{(m_2)^{(7)}((S_1)^{(7)} - (a_{38}')^{(7)})}\left[e^{(S_1)^{(7)} t} - e^{-(a_{38}')^{(7)} t}\right] + G_{38}^0 e^{-(a_{38}')^{(7)} t}\right)$$

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$$\boxed{T_{36}^0 e^{(R_1)^{(7)} t} \le T_{36}(t) \le T_{36}^0 e^{((R_1)^{(7)} + (r_{36})^{(7)})t}}$$

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$$\frac{1}{(\mu_1)^{(7)}} T_{36}^0 e^{(R_1)^{(7)} t} \le T_{36}(t) \le \frac{1}{(\mu_2)^{(7)}} T_{36}^0 e^{((R_1)^{(7)} + (r_{36})^{(7)})t}$$

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$$\frac{(b_{38})^{(7)} T_{36}^0}{(\mu_1)^{(7)}((R_1)^{(7)} - (b_{38}')^{(7)})}\left[e^{(R_1)^{(7)} t} - e^{-(b_{38}')^{(7)} t}\right] + T_{38}^0 e^{-(b_{38}')^{(7)} t} \le T_{38}(t) \le$$

$$\frac{(a_{38})^{(7)} T_{36}^0}{(\mu_2)^{(7)}((R_1)^{(7)} + (r_{36})^{(7)} + (R_2)^{(7)})}\left[e^{((R_1)^{(7)} + (r_{36})^{(7)})t} - e^{-(R_2)^{(7)} t}\right] + T_{38}^0 e^{-(R_2)^{(7)} t}$$

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| **Definition of** $(S_1)^{(7)}, (S_2)^{(7)}, (R_1)^{(7)}, (R_2)^{(7)}$:-<br><br>Where $(S_1)^{(7)} = (a_{36})^{(7)}(m_2)^{(7)} - (a'_{36})^{(7)}$<br><br>$(S_2)^{(7)} = (a_{38})^{(7)} - (p_{38})^{(7)}$<br><br>$(R_1)^{(7)} = (b_{36})^{(7)}(\mu_2)^{(7)} - (b'_{36})^{(7)}$<br><br>$(R_2)^{(7)} = (b'_{38})^{(7)} - (r_{38})^{(7)}$ | 370 |
| **Behavior of the solutions of equation**<br><br>**Theorem 2:** If we denote and define<br><br>**Definition of** $(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ :<br><br>$(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ four constants satisfying<br><br>$-(\sigma_2)^{(8)} \le -(a'_{40})^{(8)} + (a'_{41})^{(8)} - (a''_{40})^{(8)}(T_{41}, t) + (a''_{41})^{(8)}(T_{41}, t) \le -(\sigma_1)^{(8)}$<br><br>$-(\tau_2)^{(8)} \le -(b'_{40})^{(8)} + (b'_{41})^{(8)} - (b''_{40})^{(8)}\big((G_{43}), t\big) - (b''_{41})^{(8)}\big((G_{43}), t\big) \le -(\tau_1)^{(8)}$ | 371 |
| **Definition of** $(v_1)^{(8)}, (v_2)^{(8)}, (u_1)^{(8)}, (u_2)^{(8)}, v^{(8)}, u^{(8)}$ :<br><br>By $(v_1)^{(8)} > 0, (v_2)^{(8)} < 0$ and respectively $(u_1)^{(8)} > 0, (u_2)^{(8)} < 0$ the roots of the equations $(a_{41})^{(8)}\big(v^{(8)}\big)^2 + (\sigma_1)^{(8)}v^{(8)} - (a_{40})^{(8)} = 0$<br>and $(b_{41})^{(8)}\big(u^{(8)}\big)^2 + (\tau_1)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$ and | 372 |
| **Definition of** $(\bar{v}_1)^{(8)}, , (\bar{v}_2)^{(8)}, (\bar{u}_1)^{(8)}, (\bar{u}_2)^{(8)}$ :<br><br>By $(\bar{v}_1)^{(8)} > 0, (\bar{v}_2)^{(8)} < 0$ and respectively $(\bar{u}_1)^{(8)} > 0, (\bar{u}_2)^{(8)} < 0$ the<br><br>roots of the equations $(a_{41})^{(8)}\big(v^{(8)}\big)^2 + (\sigma_2)^{(8)}v^{(8)} - (a_{40})^{(8)} = 0$<br><br>and $(b_{41})^{(8)}\big(u^{(8)}\big)^2 + (\tau_2)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$<br><br>**Definition of** $(m_1)^{(8)}, (m_2)^{(8)}, (\mu_1)^{(8)}, (\mu_2)^{(8)}, (v_0)^{(8)}$ :-<br><br>If we define $(m_1)^{(8)}, (m_2)^{(8)}, (\mu_1)^{(8)}, (\mu_2)^{(8)}$ by<br><br>$(m_2)^{(8)} = (v_0)^{(8)}, (m_1)^{(8)} = (v_1)^{(8)}$, ***if*** $(v_0)^{(8)} < (v_1)^{(8)}$<br><br>$(m_2)^{(8)} = (v_1)^{(8)}, (m_1)^{(8)} = (\bar{v}_1)^{(8)}$, ***if*** $(v_1)^{(8)} < (v_0)^{(8)} < (\bar{v}_1)^{(8)}$, | |

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| and $\boxed{(v_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}}$<br><br>$(m_2)^{(8)} = (v_1)^{(8)}, (m_1)^{(8)} = (v_0)^{(8)},$ ***if*** $(\bar{v}_1)^{(8)} < (v_0)^{(8)}$ | |
| and analogously<br><br>$(\mu_2)^{(8)} = (u_0)^{(8)}, (\mu_1)^{(8)} = (u_1)^{(8)},$ ***if*** $(u_0)^{(8)} < (u_1)^{(8)}$<br><br>$(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (\bar{u}_1)^{(8)}$ , ***if*** $(u_1)^{(8)} < (u_0)^{(8)} < (\bar{u}_1)^{(8)}$,<br>and $\boxed{(u_0)^{(8)} = \frac{T_{40}^0}{T_{41}^0}}$<br><br>$(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (u_0)^{(8)},$ ***if*** $(\bar{u}_1)^{(8)} < (u_0)^{(8)}$ where $(u_1)^{(8)}, (\bar{u}_1)^{(8)}$ | 374 |
| Then the solution of global equations satisfies the inequalities<br><br>$G_{40}^0 e^{((S_1)^{(8)} - (p_{40})^{(8)})t} \leq G_{40}(t) \leq G_{40}^0 e^{(S_1)^{(8)} t}$<br><br>where $(p_i)^{(8)}$ is defined by equation | 375 |
| $\frac{1}{(m_1)^{(8)}} G_{40}^0 e^{((S_1)^{(8)} - (p_{40})^{(8)})t} \leq G_{41}(t) \leq \frac{1}{(m_2)^{(8)}} G_{40}^0 e^{(S_1)^{(8)} t}$ | 376 |
| $\left(\frac{(a_{42})^{(8)} G_{40}^0}{(m_1)^{(8)}((S_1)^{(8)} - (p_{40})^{(8)} - (S_2)^{(8)})}\left[e^{((S_1)^{(8)} - (p_{40})^{(8)})t} - e^{-(S_2)^{(8)} t}\right] + G_{42}^0 e^{-(S_2)^{(8)} t} \leq G_{42}(t) \leq \right.$<br><br>$\left.\frac{(a_{42})^{(8)} G_{40}^0}{(m_2)^{(8)}((S_1)^{(8)} - (a'_{42})^{(8)})}\left[e^{(S_1)^{(8)} t} - e^{-(a'_{42})^{(8)} t}\right] + G_{42}^0 e^{-(a'_{42})^{(8)} t}\right)$ | 377 |
| $\boxed{T_{40}^0 e^{(R_1)^{(8)} t} \leq T_{40}(t) \leq T_{40}^0 e^{((R_1)^{(8)} + (r_{40})^{(8)})t}}$ | 378 |
| $\frac{1}{(\mu_1)^{(8)}} T_{40}^0 e^{(R_1)^{(8)} t} \leq T_{40}(t) \leq \frac{1}{(\mu_2)^{(8)}} T_{40}^0 e^{((R_1)^{(8)} + (r_{40})^{(8)})t}$ | 379 |
| $\frac{(b_{42})^{(8)} T_{40}^0}{(\mu_1)^{(8)}((R_1)^{(8)} - (b'_{42})^{(8)})}\left[e^{(R_1)^{(8)} t} - e^{-(b'_{42})^{(8)} t}\right] + T_{42}^0 e^{-(b'_{42})^{(8)} t} \leq T_{42}(t) \leq$<br><br>$\frac{(a_{42})^{(8)} T_{40}^0}{(\mu_2)^{(8)}((R_1)^{(8)} + (r_{40})^{(8)} + (R_2)^{(8)})}\left[e^{((R_1)^{(8)} + (r_{40})^{(8)})t} - e^{-(R_2)^{(8)} t}\right] + T_{42}^0 e^{-(R_2)^{(8)} t}$ | 380 |
| **Definition of** $(S_1)^{(8)}, (S_2)^{(8)}, (R_1)^{(8)}, (R_2)^{(8)}$:-<br><br>Where $(S_1)^{(8)} = (a_{40})^{(8)}(m_2)^{(8)} - (a'_{40})^{(8)}$<br><br>$(S_2)^{(8)} = (a_{42})^{(8)} - (p_{42})^{(8)}$ | 381 |

$(R_1)^{(8)} = (b_{40})^{(8)}(\mu_2)^{(8)} - (b'_{40})^{(8)}$

$(R_2)^{(8)} = (b'_{42})^{(8)} - (r_{42})^{(8)}$

**Behavior of the solutions of equation 37 to 92** 382

**Theorem 2:** If we denote and define

**Definition of** $(\sigma_1)^{(9)}, (\sigma_2)^{(9)}, (\tau_1)^{(9)}, (\tau_2)^{(9)}$ :

$(\sigma_1)^{(9)}, (\sigma_2)^{(9)}, (\tau_1)^{(9)}, (\tau_2)^{(9)}$ four constants satisfying

$-(\sigma_2)^{(9)} \le -(a'_{44})^{(9)} + (a'_{45})^{(9)} - (a''_{44})^{(9)}(T_{45}, t) + (a''_{45})^{(9)}(T_{45}, t) \le -(\sigma_1)^{(9)}$

$-(\tau_2)^{(9)} \le -(b'_{44})^{(9)} + (b'_{45})^{(9)} - (b''_{44})^{(9)}\big((G_{47}), t\big) - (b''_{45})^{(9)}\big((G_{47}), t\big) \le -(\tau_1)^{(9)}$

**Definition of** $(\nu_1)^{(9)}, (\nu_2)^{(9)}, (u_1)^{(9)}, (u_2)^{(9)}, \nu^{(9)}, u^{(9)}$ :

By $(\nu_1)^{(9)} > 0$ , $(\nu_2)^{(9)} < 0$ and respectively $(u_1)^{(9)} > 0$ , $(u_2)^{(9)} < 0$ the roots of the equations $(a_{45})^{(9)}\big(\nu^{(9)}\big)^2 + (\sigma_1)^{(9)}\nu^{(9)} - (a_{44})^{(9)} = 0$ and $(b_{45})^{(9)}\big(u^{(9)}\big)^2 + (\tau_1)^{(9)}u^{(9)} - (b_{44})^{(9)} = 0$ and

**Definition of** $(\bar{\nu}_1)^{(9)}, , (\bar{\nu}_2)^{(9)}, (\bar{u}_1)^{(9)}, (\bar{u}_2)^{(9)}$ :

By $(\bar{\nu}_1)^{(9)} > 0$ , $(\bar{\nu}_2)^{(9)} < 0$ and respectively $(\bar{u}_1)^{(9)} > 0$ , $(\bar{u}_2)^{(9)} < 0$ the roots of the equations $(a_{45})^{(9)}\big(\nu^{(9)}\big)^2 + (\sigma_2)^{(9)}\nu^{(9)} - (a_{44})^{(9)} = 0$ and $(b_{45})^{(9)}\big(u^{(9)}\big)^2 + (\tau_2)^{(9)}u^{(9)} - (b_{44})^{(9)} = 0$

**Definition of** $(m_1)^{(9)}, (m_2)^{(9)}, (\mu_1)^{(9)}, (\mu_2)^{(9)}, (\nu_0)^{(9)}$ :-

If we define $(m_1)^{(9)}, (m_2)^{(9)}, (\mu_1)^{(9)}, (\mu_2)^{(9)}$ by

$(m_2)^{(9)} = (\nu_0)^{(9)}, (m_1)^{(9)} = (\nu_1)^{(9)}$, ***if*** $(\nu_0)^{(9)} < (\nu_1)^{(9)}$

$(m_2)^{(9)} = (\nu_1)^{(9)}, (m_1)^{(9)} = (\bar{\nu}_1)^{(9)}$ , ***if*** $(\nu_1)^{(9)} < (\nu_0)^{(9)} < (\bar{\nu}_1)^{(9)}$,

and $\boxed{(\nu_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}}}$

$(m_2)^{(9)} = (\nu_1)^{(9)}, (m_1)^{(9)} = (\nu_0)^{(9)}$, ***if*** $(\bar{\nu}_1)^{(9)} < (\nu_0)^{(9)}$

and analogously

$(\mu_2)^{(9)} = (u_0)^{(9)}, (\mu_1)^{(9)} = (u_1)^{(9)}$, ***if*** $(u_0)^{(9)} < (u_1)^{(9)}$

$(\mu_2)^{(9)} = (u_1)^{(9)}, (\mu_1)^{(9)} = (\bar{u}_1)^{(9)}$ , ***if*** $(u_1)^{(9)} < (u_0)^{(9)} < (\bar{u}_1)^{(9)}$,

and $\boxed{(u_0)^{(9)} = \frac{T^0_{44}}{T^0_{45}}}$

$(\mu_2)^{(9)} = (u_1)^{(9)}, (\mu_1)^{(9)} = (u_0)^{(9)}$, ***if*** $(\bar{u}_1)^{(9)} < (u_0)^{(9)}$ where $(u_1)^{(9)}, (\bar{u}_1)^{(9)}$ are defined by 59 and 69 respectively

Then the solution of 19,20,21,22,23 and 24 satisfies the inequalities

$$G_{44}^{0}e^{((S_1)^{(9)}-(p_{44})^{(9)})t} \le G_{44}(t) \le G_{44}^{0}e^{(S_1)^{(9)}t}$$

where $(p_i)^{(9)}$ is defined by equation 45

$$\frac{1}{(m_9)^{(9)}}G_{44}^{0}e^{((S_1)^{(9)}-(p_{44})^{(9)})t} \le G_{45}(t) \le \frac{1}{(m_2)^{(9)}}G_{44}^{0}e^{(S_1)^{(9)}t}$$

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$$\frac{(a_{46})^{(9)}G_{44}^{0}}{(m_1)^{(9)}((S_1)^{(9)}-(p_{44})^{(9)}-(S_2)^{(9)})}\left[e^{((S_1)^{(9)}-(p_{44})^{(9)})t}-e^{-(S_2)^{(9)}t}\right]+G_{46}^{0}e^{-(S_2)^{(9)}t} \le G_{46}(t) \le$$

$$\frac{(a_{46})^{(9)}G_{44}^{0}}{(m_2)^{(9)}((S_1)^{(9)}-(a_{46}^{'})^{(9)})}[e^{(S_1)^{(9)}t}-e^{-(a_{46}^{'})^{(9)}t}]+\ G_{46}^{0}e^{-(a_{46}^{'})^{(9)}t})$$

$$\boxed{T_{44}^{0}e^{(R_1)^{(9)}t} \le T_{44}(t) \le T_{44}^{0}e^{((R_1)^{(9)}+(r_{44})^{(9)})t}}$$

$$\frac{1}{(\mu_1)^{(9)}}T_{44}^{0}e^{(R_1)^{(9)}t} \le T_{44}(t) \le \frac{1}{(\mu_2)^{(9)}}T_{44}^{0}e^{((R_1)^{(9)}+(r_{44})^{(9)})t}$$

$$\frac{(b_{46})^{(9)}T_{44}^{0}}{(\mu_1)^{(9)}((R_1)^{(9)}-(b_{46}^{'})^{(9)})}\left[e^{(R_1)^{(9)}t}-e^{-(b_{46}^{'})^{(9)}t}\right]+T_{46}^{0}e^{-(b_{46}^{'})^{(9)}t} \le T_{46}(t) \le$$

$$\frac{(a_{46})^{(9)}T_{44}^{0}}{(\mu_2)^{(9)}((R_1)^{(9)}+(r_{44})^{(9)}+(R_2)^{(9)})}\left[e^{((R_1)^{(9)}+(r_{44})^{(9)})t}-e^{-(R_2)^{(9)}t}\right]+T_{46}^{0}e^{-(R_2)^{(9)}t}$$

**Definition of** $(S_1)^{(9)}, (S_2)^{(9)}, (R_1)^{(9)}, (R_2)^{(9)}$:-

Where $(S_1)^{(9)} = (a_{44})^{(9)}(m_2)^{(9)} - (a_{44}^{'})^{(9)}$

$(S_2)^{(9)} = (a_{46})^{(9)} - (p_{46})^{(9)}$

$(R_1)^{(9)} = (b_{44})^{(9)}(\mu_2)^{(9)} - (b_{44}^{'})^{(9)}$

$(R_2)^{(9)} = (b_{46}^{'})^{(9)} - (r_{46})^{(9)}$

**Proof :** From global equations we obtain

$$\frac{dv^{(1)}}{dt} = (a_{13})^{(1)} - \left((a_{13}^{'})^{(1)} - (a_{14}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14},t)\right) - (a_{14}^{''})^{(1)}(T_{14},t)v^{(1)} - (a_{14})^{(1)}v^{(1)}$$

**Definition of** $v^{(1)}$ :- $\boxed{v^{(1)} = \frac{G_{13}}{G_{14}}}$

It follows

$$-\left((a_{14})^{(1)}(v^{(1)})^{2} + (\sigma_2)^{(1)}v^{(1)} - (a_{13})^{(1)}\right) \le \frac{dv^{(1)}}{dt} \le -\left((a_{14})^{(1)}(v^{(1)})^{2} + (\sigma_1)^{(1)}v^{(1)} - (a_{13})^{(1)}\right)$$

From which one obtains

**Definition of** $(\bar{v}_1)^{(1)}, (v_0)^{(1)}$ :-

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For $0 < \boxed{(\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}} < (\nu_1)^{(1)} < (\bar{\nu}_1)^{(1)}$

$$\nu^{(1)}(t) \geq \frac{(\nu_1)^{(1)} + (C)^{(1)}(\nu_2)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)} - (\nu_0)^{(1)}\right)t\right]}}{1 + (C)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)} - (\nu_0)^{(1)}\right)t\right]}} \quad , \quad \boxed{(C)^{(1)} = \frac{(\nu_1)^{(1)} - (\nu_0)^{(1)}}{(\nu_0)^{(1)} - (\nu_2)^{(1)}}}$$

it follows $(\nu_0)^{(1)} \leq \nu^{(1)}(t) \leq (\nu_1)^{(1)}$

In the same manner , we get 384

$$\nu^{(1)}(t) \leq \frac{(\bar{\nu}_1)^{(1)} + (\bar{C})^{(1)}(\bar{\nu}_2)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}}{1 + (\bar{C})^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(1)} = \frac{(\bar{\nu}_1)^{(1)} - (\nu_0)^{(1)}}{(\nu_0)^{(1)} - (\bar{\nu}_2)^{(1)}}}$$

From which we deduce $(\nu_0)^{(1)} \leq \nu^{(1)}(t) \leq (\bar{\nu}_1)^{(1)}$

If $0 < (\nu_1)^{(1)} < (\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0} < (\bar{\nu}_1)^{(1)}$ we find like in the previous case, 385

$$(\nu_1)^{(1)} \leq \frac{(\nu_1)^{(1)} + (C)^{(1)}(\nu_2)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)} - (\nu_2)^{(1)}\right)t\right]}}{1 + (C)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)} - (\nu_2)^{(1)}\right)t\right]}} \leq \nu^{(1)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(1)} + (\bar{C})^{(1)}(\bar{\nu}_2)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}}{1 + (\bar{C})^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}} \leq (\bar{\nu}_1)^{(1)}$$

If $0 < (\nu_1)^{(1)} \leq (\bar{\nu}_1)^{(1)} \leq \boxed{(\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}}$ , we obtain 386

$$(\nu_1)^{(1)} \leq \nu^{(1)}(t) \leq \frac{(\bar{\nu}_1)^{(1)} + (C)^{(1)}(\bar{\nu}_2)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}}{1 + (C)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}} \leq (\nu_0)^{(1)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(1)}(t)$ :-

$$(m_2)^{(1)} \leq \nu^{(1)}(t) \leq (m_1)^{(1)}, \quad \boxed{\nu^{(1)}(t) = \frac{G_{13}(t)}{G_{14}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(1)}(t)$ :-

$$(\mu_2)^{(1)} \leq u^{(1)}(t) \leq (\mu_1)^{(1)}, \quad \boxed{u^{(1)}(t) = \frac{T_{13}(t)}{T_{14}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{13}^{''})^{(1)} = (a_{14}^{''})^{(1)}, then\ (\sigma_1)^{(1)} = (\sigma_2)^{(1)}$ and in this case $(\nu_1)^{(1)} = (\bar{\nu}_1)^{(1)}$ if in addition $(\nu_0)^{(1)} = (\nu_1)^{(1)}$ then $\nu^{(1)}(t) = (\nu_0)^{(1)}$ and as a consequence $G_{13}(t) = (\nu_0)^{(1)} G_{14}(t)$ this also defines $(\nu_0)^{(1)}$ for the special case

Analogously if $(b_{13}^{''})^{(1)} = (b_{14}^{''})^{(1)}, then\ (\tau_1)^{(1)} = (\tau_2)^{(1)}$ and then

$(u_1)^{(1)} = (\bar{u}_1)^{(1)}$ if in addition $(u_0)^{(1)} = (u_1)^{(1)}$ then $T_{13}(t) = (u_0)^{(1)} T_{14}(t)$ This is an important consequence of the relation between $(\nu_1)^{(1)}$ and $(\bar{\nu}_1)^{(1)}$, and definition of $(u_0)^{(1)}$.

**Proof** : From global equations we obtain 387

$$\frac{d\nu^{(2)}}{dt} = (a_{16})^{(2)} - \left((a_{16}^{'})^{(2)} - (a_{17}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17}, t)\right) - (a_{17}^{''})^{(2)}(T_{17}, t)\nu^{(2)} - (a_{17})^{(2)}\nu^{(2)}$$

**Definition of** $\nu^{(2)}$ :- $\boxed{\nu^{(2)} = \frac{G_{16}}{G_{17}}}$ 388

It follows 389

$$-\left((a_{17})^{(2)}\left(\nu^{(2)}\right)^2 + (\sigma_2)^{(2)}\nu^{(2)} - (a_{16})^{(2)}\right) \le \frac{d\nu^{(2)}}{dt} \le -\left((a_{17})^{(2)}\left(\nu^{(2)}\right)^2 + (\sigma_1)^{(2)}\nu^{(2)} - (a_{16})^{(2)}\right)$$

From which one obtains 390

**Definition of** $(\bar{\nu}_1)^{(2)}, (\nu_0)^{(2)}$ :-

For $0 < (\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0} < (\nu_1)^{(2)} < (\bar{\nu}_1)^{(2)}$

$$\nu^{(2)}(t) \ge \frac{(\nu_1)^{(2)} + (C)^{(2)}(\nu_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_0)^{(2)}\right)t\right]}}{1 + (C)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_0)^{(2)}\right)t\right]}} \quad , \quad \boxed{(C)^{(2)} = \frac{(\nu_1)^{(2)} - (\nu_0)^{(2)}}{(\nu_0)^{(2)} - (\nu_2)^{(2)}}}$$

it follows $(\nu_0)^{(2)} \le \nu^{(2)}(t) \le (\nu_1)^{(2)}$

In the same manner , we get 391

$$\nu^{(2)}(t) \le \frac{(\bar{\nu}_1)^{(2)} + (\bar{C})^{(2)}(\bar{\nu}_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}}{1 + (\bar{C})^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(2)} = \frac{(\bar{\nu}_1)^{(2)} - (\nu_0)^{(2)}}{(\nu_0)^{(2)} - (\bar{\nu}_2)^{(2)}}}$$

From which we deduce $(\nu_0)^{(2)} \le \nu^{(2)}(t) \le (\bar{\nu}_1)^{(2)}$ 392

If $0 < (\nu_1)^{(2)} < (\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0} < (\bar{\nu}_1)^{(2)}$ we find like in the previous case, 393

$$(\nu_1)^{(2)} \le \frac{(\nu_1)^{(2)} + (C)^{(2)}(\nu_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_2)^{(2)}\right)t\right]}}{1 + (C)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_2)^{(2)}\right)t\right]}} \le \nu^{(2)}(t) \le$$

| | |
|---|---|
| $\frac{(\bar{\nu}_1)^{(2)}+(\bar{C})^{(2)}(\bar{\nu}_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)}-(\bar{\nu}_2)^{(2)}\right)t\right]}}{1+(\bar{C})^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)}-(\bar{\nu}_2)^{(2)}\right)t\right]}} \leq (\bar{\nu}_1)^{(2)}$ | |
| If $0 < (\nu_1)^{(2)} \leq (\bar{\nu}_1)^{(2)} \leq (\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0}$ , we obtain<br><br>$(\nu_1)^{(2)} \leq \nu^{(2)}(t) \leq \frac{(\bar{\nu}_1)^{(2)}+(\bar{C})^{(2)}(\bar{\nu}_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)}-(\bar{\nu}_2)^{(2)}\right)t\right]}}{1+(\bar{C})^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)}-(\bar{\nu}_2)^{(2)}\right)t\right]}} \leq (\nu_0)^{(2)}$<br><br>And so with the notation of the first part of condition (c) , we have | 394 |
| **Definition of** $\nu^{(2)}(t)$ :-<br><br>$(m_2)^{(2)} \leq \nu^{(2)}(t) \leq (m_1)^{(2)}$, $\boxed{\nu^{(2)}(t) = \frac{G_{16}(t)}{G_{17}(t)}}$ | 395 |
| In a completely analogous way, we obtain<br><br>**Definition of** $u^{(2)}(t)$ :-<br><br>$(\mu_2)^{(2)} \leq u^{(2)}(t) \leq (\mu_1)^{(2)}$, $\boxed{u^{(2)}(t) = \frac{T_{16}(t)}{T_{17}(t)}}$ | 396 |
| Now, using this result and replacing it in global equations we get easily the result stated in the theorem. | |
| **Particular case :**<br><br>If $(a_{16}^{''})^{(2)} = (a_{17}^{''})^{(2)}, then\ (\sigma_1)^{(2)} = (\sigma_2)^{(2)}$ and in this case $(\nu_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ if in addition $(\nu_0)^{(2)} = (\nu_1)^{(2)}$ then $\nu^{(2)}(t) = (\nu_0)^{(2)}$ and as a consequence $G_{16}(t) = (\nu_0)^{(2)}G_{17}(t)$<br><br>Analogously if $(b_{16}^{''})^{(2)} = (b_{17}^{''})^{(2)}, then\ (\tau_1)^{(2)} = (\tau_2)^{(2)}$ and then<br><br>$(u_1)^{(2)} = (\bar{u}_1)^{(2)}$if in addition $(u_0)^{(2)} = (u_1)^{(2)}$ then $T_{16}(t) = (u_0)^{(2)}T_{17}(t)$ This is an important consequence of the relation between $(\nu_1)^{(2)}$ and $(\bar{\nu}_1)^{(2)}$ | 397 |
| **Proof** : From global equations we obtain<br><br>$\frac{d\nu^{(3)}}{dt} = (a_{20})^{(3)} - \left((a_{20}^{'})^{(3)} - (a_{21}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21},t)\right) - (a_{21}^{''})^{(3)}(T_{21},t)\nu^{(3)} - (a_{21})^{(3)}\nu^{(3)}$ | 398 |
| **Definition of** $\nu^{(3)}$ :- $\boxed{\nu^{(3)} = \frac{G_{20}}{G_{21}}}$<br><br>It follows<br><br>$-\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_2)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right) \leq \frac{d\nu^{(3)}}{dt} \leq -\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_1)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right)$ | 399 |
| | 400 |

From which one obtains

For $0 < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\nu_1)^{(3)} < (\bar{\nu}_1)^{(3)}$

$$\nu^{(3)}(t) \geq \frac{(\nu_1)^{(3)}+(C)^{(3)}(\nu_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_0)^{(3)}\right)t\right]}}{1+(C)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_0)^{(3)}\right)t\right]}} \quad , \quad \boxed{(C)^{(3)} = \frac{(\nu_1)^{(3)}-(\nu_0)^{(3)}}{(\nu_0)^{(3)}-(\nu_2)^{(3)}}}$$

it follows $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\nu_1)^{(3)}$

In the same manner , we get

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$$\nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(3)} = \frac{(\bar{\nu}_1)^{(3)}-(\nu_0)^{(3)}}{(\nu_0)^{(3)}-(\bar{\nu}_2)^{(3)}}}$$

**Definition of** $(\bar{\nu}_1)^{(3)}$ :-

From which we deduce $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\bar{\nu}_1)^{(3)}$

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If $0 < (\nu_1)^{(3)} < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\bar{\nu}_1)^{(3)}$ we find like in the previous case,

$$(\nu_1)^{(3)} \leq \frac{(\nu_1)^{(3)}+(C)^{(3)}(\nu_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}}{1+(C)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}} \leq \nu^{(3)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\bar{\nu}_1)^{(3)}$$

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If $0 < (\nu_1)^{(3)} \leq (\bar{\nu}_1)^{(3)} \leq (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0}$ , we obtain

$$(\nu_1)^{(3)} \leq \nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\nu_0)^{(3)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(3)}(t)$ :-

$$(m_2)^{(3)} \leq \nu^{(3)}(t) \leq (m_1)^{(3)}, \quad \boxed{\nu^{(3)}(t) = \frac{G_{20}(t)}{G_{21}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(3)}(t)$ :-

$$(\mu_2)^{(3)} \leq u^{(3)}(t) \leq (\mu_1)^{(3)}, \quad \boxed{u^{(3)}(t) = \frac{T_{20}(t)}{T_{21}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{20}^{''})^{(3)} = (a_{21}^{''})^{(3)}, then\ (\sigma_1)^{(3)} = (\sigma_2)^{(3)}$ and in this case $(\nu_1)^{(3)} = (\bar{\nu}_1)^{(3)}$ if in addition $(\nu_0)^{(3)} = (\nu_1)^{(3)}$ then $\nu^{(3)}(t) = (\nu_0)^{(3)}$ and as a consequence $G_{20}(t) = (\nu_0)^{(3)} G_{21}(t)$

Analogously if $(b_{20}^{''})^{(3)} = (b_{21}^{''})^{(3)}, then\ (\tau_1)^{(3)} = (\tau_2)^{(3)}$ and then

$(u_1)^{(3)} = (\bar{u}_1)^{(3)}$ if in addition $(u_0)^{(3)} = (u_1)^{(3)}$ then $T_{20}(t) = (u_0)^{(3)} T_{21}(t)$ This is an important consequence of the relation between $(\nu_1)^{(3)}$ and $(\bar{\nu}_1)^{(3)}$

**Proof** : From global equations we obtain

$$\frac{d\nu^{(4)}}{dt} = (a_{24})^{(4)} - \left((a_{24}^{'})^{(4)} - (a_{25}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25}, t)\right) - (a_{25}^{''})^{(4)}(T_{25}, t)\nu^{(4)} - (a_{25})^{(4)}\nu^{(4)}$$

**Definition of** $\nu^{(4)}$ :- $\boxed{\nu^{(4)} = \frac{G_{24}}{G_{25}}}$

It follows

$$-\left((a_{25})^{(4)}\left(\nu^{(4)}\right)^2 + (\sigma_2)^{(4)}\nu^{(4)} - (a_{24})^{(4)}\right) \leq \frac{d\nu^{(4)}}{dt} \leq -\left((a_{25})^{(4)}\left(\nu^{(4)}\right)^2 + (\sigma_4)^{(4)}\nu^{(4)} - (a_{24})^{(4)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(4)}, (\nu_0)^{(4)}$ :-

For $0 < \boxed{(\nu_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0}} < (\nu_1)^{(4)} < (\bar{\nu}_1)^{(4)}$

$$\nu^{(4)}(t) \geq \frac{(\nu_1)^{(4)} + (C)^{(4)}(\nu_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_0)^{(4)}\right)t\right]}}{4 + (C)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_0)^{(4)}\right)t\right]}} \quad , \quad \boxed{(C)^{(4)} = \frac{(\nu_1)^{(4)} - (\nu_0)^{(4)}}{(\nu_0)^{(4)} - (\nu_2)^{(4)}}}$$

it follows $(\nu_0)^{(4)} \leq \nu^{(4)}(t) \leq (\nu_1)^{(4)}$

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In the same manner , we get

$$\nu^{(4)}(t) \leq \frac{(\bar{\nu}_1)^{(4)} + (\bar{C})^{(4)}(\bar{\nu}_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}}{4 + (\bar{C})^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(4)} = \frac{(\bar{\nu}_1)^{(4)} - (\nu_0)^{(4)}}{(\nu_0)^{(4)} - (\bar{\nu}_2)^{(4)}}}$$

From which we deduce $(\nu_0)^{(4)} \leq \nu^{(4)}(t) \leq (\bar{\nu}_1)^{(4)}$

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If $0 < (\nu_1)^{(4)} < (\nu_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0} < (\bar{\nu}_1)^{(4)}$ we find like in the previous case,

$$(\nu_1)^{(4)} \leq \frac{(\nu_1)^{(4)} + (C)^{(4)}(\nu_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_2)^{(4)}\right)t\right]}}{1 + (C)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_2)^{(4)}\right)t\right]}} \leq \nu^{(4)}(t) \leq$$

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$$\frac{(\bar{\nu}_1)^{(4)}+(\bar{C})^{(4)}(\bar{\nu}_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)}-(\bar{\nu}_2)^{(4)}\right)t\right]}}{1+(\bar{C})^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)}-(\bar{\nu}_2)^{(4)}\right)t\right]}} \leq (\bar{\nu}_1)^{(4)}$$

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If $0 < (\nu_1)^{(4)} \leq (\bar{\nu}_1)^{(4)} \leq \boxed{(\nu_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0}}$ , we obtain

$$(\nu_1)^{(4)} \leq \nu^{(4)}(t) \leq \frac{(\bar{\nu}_1)^{(4)}+(\bar{C})^{(4)}(\bar{\nu}_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)}-(\bar{\nu}_2)^{(4)}\right)t\right]}}{1+(\bar{C})^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)}-(\bar{\nu}_2)^{(4)}\right)t\right]}} \leq (\nu_0)^{(4)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(4)}(t)$ :-

$(m_2)^{(4)} \leq \nu^{(4)}(t) \leq (m_1)^{(4)}$, $\boxed{\nu^{(4)}(t) = \frac{G_{24}(t)}{G_{25}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(4)}(t)$ :-

$(\mu_2)^{(4)} \leq u^{(4)}(t) \leq (\mu_1)^{(4)}$, $\boxed{u^{(4)}(t) = \frac{T_{24}(t)}{T_{25}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{24}^{''})^{(4)} = (a_{25}^{''})^{(4)}, then\ (\sigma_1)^{(4)} = (\sigma_2)^{(4)}$ and in this case $(\nu_1)^{(4)} = (\bar{\nu}_1)^{(4)}$ if in addition $(\nu_0)^{(4)} = (\nu_1)^{(4)}$ then $\nu^{(4)}(t) = (\nu_0)^{(4)}$ and as a consequence $G_{24}(t) = (\nu_0)^{(4)}G_{25}(t)$ **this also defines** $(\nu_0)^{(4)}$ **for the special case** .

Analogously if $(b_{24}^{''})^{(4)} = (b_{25}^{''})^{(4)}, then\ (\tau_1)^{(4)} = (\tau_2)^{(4)}$ and then $(u_1)^{(4)} = (\bar{u}_4)^{(4)}$ if in addition $(u_0)^{(4)} = (u_1)^{(4)}$ then $T_{24}(t) = (u_0)^{(4)}T_{25}(t)$ This is an important consequence of the relation between $(\nu_1)^{(4)}$ and $(\bar{\nu}_1)^{(4)}$, **and definition of** $(u_0)^{(4)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(5)}}{dt} = (a_{28})^{(5)} - \left((a_{28}^{'})^{(5)} - (a_{29}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}, t)\right) - (a_{29}^{''})^{(5)}(T_{29}, t)\nu^{(5)} - (a_{29})^{(5)}\nu^{(5)}$$

**Definition of** $\nu^{(5)}$ :- $\boxed{\nu^{(5)} = \frac{G_{28}}{G_{29}}}$

It follows

$$-\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_2)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right) \leq \frac{d\nu^{(5)}}{dt} \leq -\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_1)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(5)}, (\nu_0)^{(5)}$ :-

For $0 < \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}} < (\nu_1)^{(5)} < (\bar{\nu}_1)^{(5)}$

$$\nu^{(5)}(t) \geq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}}{5 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}} \quad , \quad \boxed{(C)^{(5)} = \frac{(\nu_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\nu_2)^{(5)}}}$$

it follows $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\nu_1)^{(5)}$

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In the same manner , we get

$$\nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)} + (\bar{C})^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{5 + (\bar{C})^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(5)} = \frac{(\bar{\nu}_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\bar{\nu}_2)^{(5)}}}$$

From which we deduce $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\bar{\nu}_5)^{(5)}$

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If $0 < (\nu_1)^{(5)} < (\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0} < (\bar{\nu}_1)^{(5)}$ we find like in the previous case,

$$(\nu_1)^{(5)} \leq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}}{1 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}} \leq \nu^{(5)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(5)} + (C)^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{1 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\bar{\nu}_1)^{(5)}$$

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If $0 < (\nu_1)^{(5)} \leq (\bar{\nu}_1)^{(5)} \leq \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}}$ , we obtain

$$(\nu_1)^{(5)} \leq \nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)} + (C)^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{1 + (\bar{C})^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\nu_0)^{(5)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(5)}(t)$ :-

$$(m_2)^{(5)} \leq \nu^{(5)}(t) \leq (m_1)^{(5)}, \quad \boxed{\nu^{(5)}(t) = \frac{G_{28}(t)}{G_{29}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(5)}(t)$ :-

$$(\mu_2)^{(5)} \leq u^{(5)}(t) \leq (\mu_1)^{(5)}, \quad \boxed{u^{(5)}(t) = \frac{T_{28}(t)}{T_{29}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{28}'')^{(5)} = (a_{29}'')^{(5)}, then\ (\sigma_1)^{(5)} = (\sigma_2)^{(5)}$ and in this case $(\nu_1)^{(5)} = (\bar{\nu}_1)^{(5)}$ if in addition $(\nu_0)^{(5)} = (\nu_5)^{(5)}$ then $\nu^{(5)}(t) = (\nu_0)^{(5)}$ and as a consequence $G_{28}(t) = (\nu_0)^{(5)} G_{29}(t)$ **this also defines** $(\nu_0)^{(5)}$ **for the special case** .

Analogously if $(b_{28}'')^{(5)} = (b_{29}'')^{(5)}, then\ (\tau_1)^{(5)} = (\tau_2)^{(5)}$ and then

$(u_1)^{(5)} = (\bar{u}_1)^{(5)}$if in addition $(u_0)^{(5)} = (u_1)^{(5)}$ then $T_{28}(t) = (u_0)^{(5)}T_{29}(t)$ This is an important consequence of the relation between $(v_1)^{(5)}$ and $(\bar{v}_1)^{(5)}$, **and definition of** $(u_0)^{(5)}$.

**Proof** : From global equations we obtain 412

$$\frac{dv^{(6)}}{dt} = (a_{32})^{(6)} - \left((a'_{32})^{(6)} - (a'_{33})^{(6)} + (a''_{32})^{(6)}(T_{33},t)\right) - (a''_{33})^{(6)}(T_{33},t)v^{(6)} - (a_{33})^{(6)}v^{(6)}$$

**Definition of** $v^{(6)}$ :- $\boxed{v^{(6)} = \frac{G_{32}}{G_{33}}}$

It follows

$$-\left((a_{33})^{(6)}\left(v^{(6)}\right)^2 + (\sigma_2)^{(6)}v^{(6)} - (a_{32})^{(6)}\right) \le \frac{dv^{(6)}}{dt} \le -\left((a_{33})^{(6)}\left(v^{(6)}\right)^2 + (\sigma_1)^{(6)}v^{(6)} - (a_{32})^{(6)}\right)$$

From which one obtains

**Definition of** $(\bar{v}_1)^{(6)}, (v_0)^{(6)}$ :-

For $0 < \boxed{(v_0)^{(6)} = \frac{G^0_{32}}{G^0_{33}}} < (v_1)^{(6)} < (\bar{v}_1)^{(6)}$

$$v^{(6)}(t) \ge \frac{(v_1)^{(6)} + (C)^{(6)}(v_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((v_1)^{(6)} - (v_0)^{(6)}\right)t\right]}}{1 + (C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((v_1)^{(6)} - (v_0)^{(6)}\right)t\right]}} \quad , \quad \boxed{(C)^{(6)} = \frac{(v_1)^{(6)} - (v_0)^{(6)}}{(v_0)^{(6)} - (v_2)^{(6)}}}$$

it follows $(v_0)^{(6)} \le v^{(6)}(t) \le (v_1)^{(6)}$

In the same manner , we get 413

$$v^{(6)}(t) \le \frac{(\bar{v}_1)^{(6)} + (\bar{C})^{(6)}(\bar{v}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{v}_1)^{(6)} - (\bar{v}_2)^{(6)}\right)t\right]}}{1 + (\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{v}_1)^{(6)} - (\bar{v}_2)^{(6)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(6)} = \frac{(\bar{v}_1)^{(6)} - (v_0)^{(6)}}{(v_0)^{(6)} - (\bar{v}_2)^{(6)}}}$$

From which we deduce $(v_0)^{(6)} \le v^{(6)}(t) \le (\bar{v}_1)^{(6)}$

If $0 < (v_1)^{(6)} < (v_0)^{(6)} = \frac{G^0_{32}}{G^0_{33}} < (\bar{v}_1)^{(6)}$ we find like in the previous case, 414

$$(v_1)^{(6)} \le \frac{(v_1)^{(6)} + (C)^{(6)}(v_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((v_1)^{(6)} - (v_2)^{(6)}\right)t\right]}}{1 + (C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((v_1)^{(6)} - (v_2)^{(6)}\right)t\right]}} \le v^{(6)}(t) \le$$

$$\frac{(\bar{v}_1)^{(6)} + (\bar{C})^{(6)}(\bar{v}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{v}_1)^{(6)} - (\bar{v}_2)^{(6)}\right)t\right]}}{1 + (\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{v}_1)^{(6)} - (\bar{v}_2)^{(6)}\right)t\right]}} \le (\bar{v}_1)^{(6)}$$

If $0 < (v_1)^{(6)} \le (\bar{v}_1)^{(6)} \le \boxed{(v_0)^{(6)} = \frac{G^0_{32}}{G^0_{33}}}$ , we obtain 415

$$(v_1)^{(6)} \le v^{(6)}(t) \le \frac{(\bar{v}_1)^{(6)} + (\bar{C})^{(6)}(\bar{v}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{v}_1)^{(6)} - (\bar{v}_2)^{(6)}\right)t\right]}}{1 + (C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{v}_1)^{(6)} - (\bar{v}_2)^{(6)}\right)t\right]}} \le (v_0)^{(6)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $v^{(6)}(t)$ :-

$(m_2)^{(6)} \leq \nu^{(6)}(t) \leq (m_1)^{(6)}$, $\boxed{\nu^{(6)}(t) = \frac{G_{32}(t)}{G_{33}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(6)}(t)$ :-

$(\mu_2)^{(6)} \leq u^{(6)}(t) \leq (\mu_1)^{(6)}$, $\boxed{u^{(6)}(t) = \frac{T_{32}(t)}{T_{33}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{32}^{''})^{(6)} = (a_{33}^{''})^{(6)}, then\ (\sigma_1)^{(6)} = (\sigma_2)^{(6)}$ and in this case $(\nu_1)^{(6)} = (\bar{\nu}_1)^{(6)}$ if in addition $(\nu_0)^{(6)} = (\nu_1)^{(6)}$ then $\nu^{(6)}(t) = (\nu_0)^{(6)}$ and as a consequence $G_{32}(t) = (\nu_0)^{(6)} G_{33}(t)$ **this also defines** $(\nu_0)^{(6)}$ **for the special case** .

Analogously if $(b_{32}^{''})^{(6)} = (b_{33}^{''})^{(6)}, then\ (\tau_1)^{(6)} = (\tau_2)^{(6)}$ and then $(u_1)^{(6)} = (\bar{u}_1)^{(6)}$ if in addition $(u_0)^{(6)} = (u_1)^{(6)}$ then $T_{32}(t) = (u_0)^{(6)} T_{33}(t)$ This is an important consequence of the relation between $(\nu_1)^{(6)}$ and $(\bar{\nu}_1)^{(6)}$, **and definition of** $(u_0)^{(6)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(7)}}{dt} = (a_{36})^{(7)} - \left((a_{36}^{'})^{(7)} - (a_{37}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37}, t)\right) - (a_{37}^{''})^{(7)}(T_{37}, t)\nu^{(7)} - (a_{37})^{(7)}\nu^{(7)}$$

**Definition of** $\nu^{(7)}$ :- $\boxed{\nu^{(7)} = \frac{G_{36}}{G_{37}}}$

It follows

$$-\left((a_{37})^{(7)}\left(\nu^{(7)}\right)^2 + (\sigma_2)^{(7)}\nu^{(7)} - (a_{36})^{(7)}\right) \leq \frac{d\nu^{(7)}}{dt} \leq -\left((a_{37})^{(7)}\left(\nu^{(7)}\right)^2 + (\sigma_1)^{(7)}\nu^{(7)} - (a_{36})^{(7)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(7)}, (\nu_0)^{(7)}$ :-

For $0 < \boxed{(\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0}} < (\nu_1)^{(7)} < (\bar{\nu}_1)^{(7)}$

$$\nu^{(7)}(t) \geq \frac{(\nu_1)^{(7)} + (C)^{(7)}(\nu_2)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)} - (\nu_0)^{(7)}\right)t\right]}}{1 + (C)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)} - (\nu_0)^{(7)}\right)t\right]}} \quad , \quad \boxed{(C)^{(7)} = \frac{(\nu_1)^{(7)} - (\nu_0)^{(7)}}{(\nu_0)^{(7)} - (\nu_2)^{(7)}}}$$

it follows $(\nu_0)^{(7)} \leq \nu^{(7)}(t) \leq (\nu_1)^{(7)}$

In the same manner , we get

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$$\nu^{(7)}(t) \le \frac{(\bar{\nu}_1)^{(7)}+(\bar{C})^{(7)}(\bar{\nu}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}}{1+(\bar{C})^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(7)} = \frac{(\bar{\nu}_1)^{(7)}-(\nu_0)^{(7)}}{(\nu_0)^{(7)}-(\bar{\nu}_2)^{(7)}}}$$

From which we deduce $(\nu_0)^{(7)} \le \nu^{(7)}(t) \le (\bar{\nu}_1)^{(7)}$

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If $0 < (\nu_1)^{(7)} < (\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0} < (\bar{\nu}_1)^{(7)}$ we find like in the previous case,

$$(\nu_1)^{(7)} \le \frac{(\nu_1)^{(7)}+(C)^{(7)}(\nu_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)}-(\nu_2)^{(7)}\right)t\right]}}{1+(C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)}-(\nu_2)^{(7)}\right)t\right]}} \le \nu^{(7)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(7)}+(\bar{C})^{(7)}(\bar{\nu}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}}{1+(\bar{C})^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}} \le (\bar{\nu}_1)^{(7)}$$

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If $0 < (\nu_1)^{(7)} \le (\bar{\nu}_1)^{(7)} \le \boxed{(\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0}}$ , we obtain

$$(\nu_1)^{(7)} \le \nu^{(7)}(t) \le \frac{(\bar{\nu}_1)^{(7)}+(C)^{(7)}(\bar{\nu}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}}{1+(C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}} \le (\nu_0)^{(7)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(7)}(t)$ :-

$$(m_2)^{(7)} \le \nu^{(7)}(t) \le (m_1)^{(7)}, \quad \boxed{\nu^{(7)}(t) = \frac{G_{36}(t)}{G_{37}(t)}}$$

In a completely analogous way, we obtain

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**Definition of** $u^{(7)}(t)$ :-

$$(\mu_2)^{(7)} \le u^{(7)}(t) \le (\mu_1)^{(7)}, \quad \boxed{u^{(7)}(t) = \frac{T_{36}(t)}{T_{37}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{36}'')^{(7)} = (a_{37}'')^{(7)}, then\ (\sigma_1)^{(7)} = (\sigma_2)^{(7)}$ and in this case $(\nu_1)^{(7)} = (\bar{\nu}_1)^{(7)}$ if in addition $(\nu_0)^{(7)} = (\nu_1)^{(7)}$ then $\nu^{(7)}(t) = (\nu_0)^{(7)}$ and as a consequence $G_{36}(t) = (\nu_0)^{(7)}G_{37}(t)$ **this also defines** $(\nu_0)^{(7)}$ **for the special case** .

Analogously if $(b_{36}'')^{(7)} = (b_{37}'')^{(7)}, then\ (\tau_1)^{(7)} = (\tau_2)^{(7)}$ and then $(u_1)^{(7)} = (\bar{u}_1)^{(7)}$if in addition $(u_0)^{(7)} = (u_1)^{(7)}$ then $T_{36}(t) = (u_0)^{(7)}T_{37}(t)$ This is an important consequence of the relation between $(\nu_1)^{(7)}$ and $(\bar{\nu}_1)^{(7)}$, **and definition of** $(u_0)^{(7)}$.

**Proof** : From global equations we obtain

$$\frac{dv^{(8)}}{dt} = (a_{40})^{(8)} - \left((a_{40}^{'})^{(8)} - (a_{41}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41},t)\right) - (a_{41}^{''})^{(8)}(T_{41},t)v^{(8)} - (a_{41})^{(8)}v^{(8)}$$

**Definition of** $v^{(8)}$ :- $\boxed{v^{(8)} = \frac{G_{40}}{G_{41}}}$

It follows

$$-\left((a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_2)^{(8)}v^{(8)} - (a_{40})^{(8)}\right) \leq \frac{dv^{(8)}}{dt} \leq -\left((a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_1)^{(8)}v^{(8)} - (a_{40})^{(8)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(8)}, (\nu_0)^{(8)}$ :-

For $0 < \boxed{(\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}} < (\nu_1)^{(8)} < (\bar{\nu}_1)^{(8)}$

$$\nu^{(8)}(t) \geq \frac{(\nu_1)^{(8)} + (C)^{(8)}(\nu_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}}{1 + (C)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}} \quad , \quad \boxed{(C)^{(8)} = \frac{(\nu_1)^{(8)} - (\nu_0)^{(8)}}{(\nu_0)^{(8)} - (\nu_2)^{(8)}}}$$

it follows $(\nu_0)^{(8)} \leq \nu^{(8)}(t) \leq (\nu_1)^{(8)}$

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In the same manner , we get

$$\nu^{(8)}(t) \leq \frac{(\bar{\nu}_1)^{(8)} + (\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}}{1 + (\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(8)} = \frac{(\bar{\nu}_1)^{(8)} - (\nu_0)^{(8)}}{(\nu_0)^{(8)} - (\bar{\nu}_2)^{(8)}}}$$

From which we deduce $(\nu_0)^{(8)} \leq \nu^{(8)}(t) \leq (\bar{\nu}_8)^{(8)}$

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If $0 < (\nu_1)^{(8)} < (\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0} < (\bar{\nu}_1)^{(8)}$ we find like in the previous case,

$$(\nu_1)^{(8)} \leq \frac{(\nu_1)^{(8)} + (C)^{(8)}(\nu_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_2)^{(8)}\right)t\right]}}{1 + (C)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_2)^{(8)}\right)t\right]}} \leq \nu^{(8)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(8)} + (C)^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}}{1 + (\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}} \leq (\bar{\nu}_1)^{(8)}$$

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If $0 < (\nu_1)^{(8)} \leq (\bar{\nu}_1)^{(8)} \leq \boxed{(\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}}$ , we obtain

$$(\nu_1)^{(8)} \leq \nu^{(8)}(t) \leq \frac{(\bar{\nu}_1)^{(8)} + (\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}}{1 + (\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}} \leq (\nu_0)^{(8)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $v^{(8)}(t)$ :-

$(m_2)^{(8)} \leq v^{(8)}(t) \leq (m_1)^{(8)}$, $\boxed{v^{(8)}(t) = \frac{G_{40}(t)}{G_{41}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(8)}(t)$ :-

$(\mu_2)^{(8)} \leq u^{(8)}(t) \leq (\mu_1)^{(8)}$, $\boxed{u^{(8)}(t) = \frac{T_{40}(t)}{T_{41}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{40}^{''})^{(8)} = (a_{41}^{''})^{(8)}, then\ (\sigma_1)^{(8)} = (\sigma_2)^{(8)}$ and in this case $(\nu_1)^{(8)} = (\bar{\nu}_1)^{(8)}$ if in addition $(\nu_0)^{(8)} = (\nu_1)^{(8)}$ then $v^{(8)}(t) = (\nu_0)^{(8)}$ and as a consequence $G_{40}(t) = (\nu_0)^{(8)} G_{41}(t)$ **this also defines** $(\nu_0)^{(8)}$ **for the special case** .

Analogously if $(b_{40}^{''})^{(8)} = (b_{41}^{''})^{(8)}, then\ (\tau_1)^{(8)} = (\tau_2)^{(8)}$ and then
$(u_1)^{(8)} = (\bar{u}_1)^{(8)}$if in addition $(u_0)^{(8)} = (u_1)^{(8)}$ then $T_{40}(t) = (u_0)^{(8)} T_{41}(t)$ This is an important consequence of the relation between $(\nu_1)^{(8)}$ and $(\bar{\nu}_1)^{(8)}$, **and definition of** $(u_0)^{(8)}$.

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**Proof** : From 99,20,44,22,23,44 we obtain

$$\frac{dv^{(9)}}{dt} = (a_{44})^{(9)} - \left((a_{44}^{'})^{(9)} - (a_{45}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45}, t)\right) - (a_{45}^{''})^{(9)}(T_{45}, t)\nu^{(9)} - (a_{45})^{(9)}\nu^{(9)}$$

**Definition of** $\nu^{(9)}$ :- $\boxed{\nu^{(9)} = \frac{G_{44}}{G_{45}}}$

It follows

$$-\left((a_{45})^{(9)}\left(\nu^{(9)}\right)^2 + (\sigma_2)^{(9)}\nu^{(9)} - (a_{44})^{(9)}\right) \leq \frac{d\nu^{(9)}}{dt} \leq -\left((a_{45})^{(9)}\left(\nu^{(9)}\right)^2 + (\sigma_1)^{(9)}\nu^{(9)} - (a_{44})^{(9)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(9)}, (\nu_0)^{(9)}$ :-

For $0 < \boxed{(\nu_0)^{(9)} = \frac{G_{44}^0}{G_{45}^0}} < (\nu_1)^{(9)} < (\bar{\nu}_1)^{(9)}$

$$v^{(9)}(t) \geq \frac{(\nu_1)^{(9)} + (C)^{(9)}(\nu_2)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_0)^{(9)}\right)t\right]}}{1 + (C)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_0)^{(9)}\right)t\right]}} \quad , \quad \boxed{(C)^{(9)} = \frac{(\nu_1)^{(9)} - (\nu_0)^{(9)}}{(\nu_0)^{(9)} - (\nu_2)^{(9)}}}$$

it follows $(\nu_0)^{(9)} \leq \nu^{(9)}(t) \leq (\nu_9)^{(9)}$

In the same manner , we get

$$v^{(9)}(t) \le \frac{(\bar{v}_1)^{(9)}+(\bar{C})^{(9)}(\bar{v}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)}-(\bar{v}_2)^{(9)}\right)t\right]}}{1+(\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)}-(\bar{v}_2)^{(9)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(9)} = \frac{(\bar{v}_1)^{(9)}-(v_0)^{(9)}}{(v_0)^{(9)}-(\bar{v}_2)^{(9)}}}$$

From which we deduce $(v_0)^{(9)} \le v^{(9)}(t) \le (\bar{v}_1)^{(9)}$

If $0 < (v_1)^{(9)} < (v_0)^{(9)} = \frac{G_{44}^0}{G_{45}^0} < (\bar{v}_1)^{(9)}$ we find like in the previous case,

$$(v_1)^{(9)} \le \frac{(v_1)^{(9)}+(C)^{(9)}(v_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((v_1)^{(9)}-(v_2)^{(9)}\right)t\right]}}{1+(C)^{(9)}e^{\left[-(a_{45})^{(9)}\left((v_1)^{(9)}-(v_2)^{(9)}\right)t\right]}} \le v^{(9)}(t) \le$$

$$\frac{(\bar{v}_1)^{(9)}+(\bar{C})^{(9)}(\bar{v}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)}-(\bar{v}_2)^{(9)}\right)t\right]}}{1+(\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)}-(\bar{v}_2)^{(9)}\right)t\right]}} \le (\bar{v}_1)^{(9)}$$

If $0 < (v_1)^{(9)} \le (\bar{v}_1)^{(9)} \le \boxed{(v_0)^{(9)} = \frac{G_{44}^0}{G_{45}^0}}$, we obtain

$$(v_1)^{(9)} \le v^{(9)}(t) \le \frac{(\bar{v}_1)^{(9)}+(\bar{C})^{(9)}(\bar{v}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)}-(\bar{v}_2)^{(9)}\right)t\right]}}{1+(\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)}-(\bar{v}_2)^{(9)}\right)t\right]}} \le (v_0)^{(9)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $v^{(9)}(t)$ :-

$$(m_2)^{(9)} \le v^{(9)}(t) \le (m_1)^{(9)}, \quad \boxed{v^{(9)}(t) = \frac{G_{44}(t)}{G_{45}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(9)}(t)$ :-

$$(\mu_2)^{(9)} \le u^{(9)}(t) \le (\mu_1)^{(9)}, \quad \boxed{u^{(9)}(t) = \frac{T_{44}(t)}{T_{45}(t)}}$$

Now, using this result and replacing it in 99, 20,44,22,23, and 44 we get easily the result stated in the theorem.

**Particular case :**

If $(a_{44}^{''})^{(9)} = (a_{45}^{''})^{(9)}, then\ (\sigma_1)^{(9)} = (\sigma_2)^{(9)}$ and in this case $(v_1)^{(9)} = (\bar{v}_1)^{(9)}$ if in addition $(v_0)^{(9)} = (v_1)^{(9)}$ then $v^{(9)}(t) = (v_0)^{(9)}$ and as a consequence $G_{44}(t) = (v_0)^{(9)}G_{45}(t)$ **this also defines** $(v_0)^{(9)}$ **for the special case** .

Analogously if $(b_{44}^{''})^{(9)} = (b_{45}^{''})^{(9)}, then\ (\tau_1)^{(9)} = (\tau_2)^{(9)}$ and then $(u_1)^{(9)} = (\bar{u}_1)^{(9)}$ if in addition $(u_0)^{(9)} = (u_1)^{(9)}$ then $T_{44}(t) = (u_0)^{(9)}T_{45}(t)$ This is an important consequence of the relation between $(v_1)^{(9)}$ and $(\bar{v}_1)^{(9)}$, **and definition of** $(u_0)^{(9)}$.

We can prove the following

**Theorem :** If $(a_i^{''})^{(1)} and\ (b_i^{''})^{(1)}$ are independent on $t$ , and the conditions with the notations

$$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} < 0$$

$$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a_{13})^{(1)}(p_{13})^{(1)} + (a_{14}^{'})^{(1)}(p_{14})^{(1)} + (p_{13})^{(1)}(p_{14})^{(1)} > 0$$

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| $(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} > 0$ ,<br>$(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} - (b_{13}^{'})^{(1)}(r_{14})^{(1)} - (b_{14}^{'})^{(1)}(r_{14})^{(1)} + (r_{13})^{(1)}(r_{14})^{(1)} < 0$<br>$with\ \ (p_{13})^{(1)}, (r_{14})^{(1)}$ as defined by equation are satisfied , then the system | |
| **Theorem :** If $(a_i^{''})^{(2)} and\ (b_i^{''})^{(2)}$ are independent on t , and the conditions with the notations | 426 |
| $(a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} < 0$ | 427 |
| $(a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a_{16})^{(2)}(p_{16})^{(2)} + (a_{17}^{'})^{(2)}(p_{17})^{(2)} + (p_{16})^{(2)}(p_{17})^{(2)} > 0$ | 428 |
| $(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} > 0$ , | 429 |
| $(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} - (b_{16}^{'})^{(2)}(r_{17})^{(2)} - (b_{17}^{'})^{(2)}(r_{17})^{(2)} + (r_{16})^{(2)}(r_{17})^{(2)} < 0$<br>$with\ \ (p_{16})^{(2)}, (r_{17})^{(2)}$ as defined by equation are satisfied , then the system | 430 |
| **Theorem :** If $(a_i^{''})^{(3)} and\ (b_i^{''})^{(3)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} < 0$<br>$(a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a_{20})^{(3)}(p_{20})^{(3)} + (a_{21}^{'})^{(3)}(p_{21})^{(3)} + (p_{20})^{(3)}(p_{21})^{(3)} > 0$<br>$(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} > 0$ ,<br>$(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} - (b_{20}^{'})^{(3)}(r_{21})^{(3)} - (b_{21}^{'})^{(3)}(r_{21})^{(3)} + (r_{20})^{(3)}(r_{21})^{(3)} < 0$<br>$with\ \ (p_{20})^{(3)}, (r_{21})^{(3)}$ as defined by equation are satisfied , then the system | 431 |
| We can prove the following<br>**Theorem :** If $(a_i^{''})^{(4)} and\ (b_i^{''})^{(4)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} < 0$<br>$(a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a_{24})^{(4)}(p_{24})^{(4)} + (a_{25}^{'})^{(4)}(p_{25})^{(4)} + (p_{24})^{(4)}(p_{25})^{(4)} > 0$<br>$(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} > 0$ ,<br>$(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} - (b_{24}^{'})^{(4)}(r_{25})^{(4)} - (b_{25}^{'})^{(4)}(r_{25})^{(4)} + (r_{24})^{(4)}(r_{25})^{(4)} < 0$<br>$with\ \ (p_{24})^{(4)}, (r_{25})^{(4)}$ as defined by equation are satisfied , then the system | 432 |
| **Theorem :** If $(a_i^{''})^{(5)} and\ (b_i^{''})^{(5)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} < 0$<br>$(a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a_{28})^{(5)}(p_{28})^{(5)} + (a_{29}^{'})^{(5)}(p_{29})^{(5)} + (p_{28})^{(5)}(p_{29})^{(5)} > 0$ | 433 |

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| $(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} > 0$ ,<br>$(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} - (b_{28}^{'})^{(5)}(r_{29})^{(5)} - (b_{29}^{'})^{(5)}(r_{29})^{(5)} + (r_{28})^{(5)}(r_{29})^{(5)} < 0$<br>$with\ \ (p_{28})^{(5)}, (r_{29})^{(5)}$ as defined by equation are satisfied , then the system | |
| **Theorem** If $(a_i^{''})^{(6)} and\ (b_i^{''})^{(6)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} < 0$<br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32})^{(6)}(p_{32})^{(6)} + (a_{33}^{'})^{(6)}(p_{33})^{(6)} + (p_{32})^{(6)}(p_{33})^{(6)} > 0$<br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} > 0$ ,<br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} - (b_{32}^{'})^{(6)}(r_{33})^{(6)} - (b_{33}^{'})^{(6)}(r_{33})^{(6)} + (r_{32})^{(6)}(r_{33})^{(6)} < 0$<br>$with\ \ (p_{32})^{(6)}, (r_{33})^{(6)}$ as defined by equation are satisfied , then the system | 434 |
| **Theorem :** If $(a_i^{''})^{(7)} and\ (b_i^{''})^{(7)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} < 0$<br>$(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36})^{(7)}(p_{36})^{(7)} + (a_{37}^{'})^{(7)}(p_{37})^{(7)} + (p_{36})^{(7)}(p_{37})^{(7)} > 0$<br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} > 0$ ,<br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} - (b_{36}^{'})^{(7)}(r_{37})^{(7)} - (b_{37}^{'})^{(7)}(r_{37})^{(7)} + (r_{36})^{(7)}(r_{37})^{(7)} < 0$<br>$with\ \ (p_{36})^{(7)}, (r_{37})^{(7)}$ as defined by equation are satisfied , then the system | 435 |
| **Theorem :** If $(a_i^{''})^{(8)} and\ (b_i^{''})^{(8)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} < 0$<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40})^{(8)}(p_{40})^{(8)} + (a_{41}^{'})^{(8)}(p_{41})^{(8)} + (p_{40})^{(8)}(p_{41})^{(8)} > 0$<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} > 0$ ,<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} - (b_{40}^{'})^{(8)}(r_{41})^{(8)} - (b_{41}^{'})^{(8)}(r_{41})^{(8)} + (r_{40})^{(8)}(r_{41})^{(8)} < 0$<br>$with\ \ (p_{40})^{(8)}, (r_{41})^{(8)}$ as defined by equation are satisfied , then the system | 436 |
| **Theorem :** If $(a_i^{''})^{(9)} and\ (b_i^{''})^{(9)}$ are independent on $t$ , and the conditions (with the notations 45,46,27,28)<br>$(a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} < 0$ | 436 A |

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| $(a'_{44})^{(9)}(a'_{45})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44})^{(9)}(p_{44})^{(9)} + (a'_{45})^{(9)}(p_{45})^{(9)} + (p_{44})^{(9)}(p_{45})^{(9)} > 0$<br>$(b'_{44})^{(9)}(b'_{45})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} > 0$ ,<br>$(b'_{44})^{(9)}(b'_{45})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} - (b'_{44})^{(9)}(r_{45})^{(9)} - (b'_{45})^{(9)}(r_{45})^{(9)} + (r_{44})^{(9)}(r_{45})^{(9)} < 0$<br>$with\ \ (p_{44})^{(9)}, (r_{45})^{(9)}$ as defined by equation 45 are satisfied , then the system | |
| $(a_{13})^{(1)}G_{14} - \left[(a'_{13})^{(1)} + (a''_{13})^{(1)}(T_{14})\right]G_{13} = 0$ | 437 |
| $(a_{14})^{(1)}G_{13} - \left[(a'_{14})^{(1)} + (a''_{14})^{(1)}(T_{14})\right]G_{14} = 0$ | 438 |
| $(a_{15})^{(1)}G_{14} - \left[(a'_{15})^{(1)} + (a''_{15})^{(1)}(T_{14})\right]G_{15} = 0$ | 439 |
| $(b_{13})^{(1)}T_{14} - [(b'_{13})^{(1)} - (b''_{13})^{(1)}(G)\ ]T_{13} = 0$ | 440 |
| $(b_{14})^{(1)}T_{13} - [(b'_{14})^{(1)} - (b''_{14})^{(1)}(G)\ ]T_{14} = 0$ | 441 |
| $(b_{15})^{(1)}T_{14} - [(b'_{15})^{(1)} - (b''_{15})^{(1)}(G)\ ]T_{15} = 0$ | 442 |
| has a unique positive solution , which is an equilibrium solution for the system | |
| $(a_{16})^{(2)}G_{17} - \left[(a'_{16})^{(2)} + (a''_{16})^{(2)}(T_{17})\right]G_{16} = 0$ | 443 |
| $(a_{17})^{(2)}G_{16} - \left[(a'_{17})^{(2)} + (a''_{17})^{(2)}(T_{17})\right]G_{17} = 0$ | 444 |
| $(a_{18})^{(2)}G_{17} - \left[(a'_{18})^{(2)} + (a''_{18})^{(2)}(T_{17})\right]G_{18} = 0$ | 445 |
| $(b_{16})^{(2)}T_{17} - [(b'_{16})^{(2)} - (b''_{16})^{(2)}(G_{19})\ ]T_{16} = 0$ | 446 |
| $(b_{17})^{(2)}T_{16} - [(b'_{17})^{(2)} - (b''_{17})^{(2)}(G_{19})\ ]T_{17} = 0$ | 447 |
| $(b_{18})^{(2)}T_{17} - [(b'_{18})^{(2)} - (b''_{18})^{(2)}(G_{19})\ ]T_{18} = 0$ | 448 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{20})^{(3)}G_{21} - \left[(a'_{20})^{(3)} + (a''_{20})^{(3)}(T_{21})\right]G_{20} = 0$ | 449 |
| $(a_{21})^{(3)}G_{20} - \left[(a'_{21})^{(3)} + (a''_{21})^{(3)}(T_{21})\right]G_{21} = 0$ | 450 |
| $(a_{22})^{(3)}G_{21} - \left[(a'_{22})^{(3)} + (a''_{22})^{(3)}(T_{21})\right]G_{22} = 0$ | 451 |
| $(b_{20})^{(3)}T_{21} - [(b'_{20})^{(3)} - (b''_{20})^{(3)}(G_{23})\ ]T_{20} = 0$ | 452 |
| $(b_{21})^{(3)}T_{20} - [(b'_{21})^{(3)} - (b''_{21})^{(3)}(G_{23})\ ]T_{21} = 0$ | 453 |
| $(b_{22})^{(3)}T_{21} - [(b'_{22})^{(3)} - (b''_{22})^{(3)}(G_{23})\ ]T_{22} = 0$ | 454 |
| has a unique positive solution , which is an equilibrium solution | |

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| $(a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25})\right]G_{24} = 0$ | 455 |
| $(a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25})\right]G_{25} = 0$ | 456 |
| $(a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25})\right]G_{26} = 0$ | 457 |
| $(b_{24})^{(4)}T_{25} - [(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}\left((G_{27})\right)]T_{24} = 0$ | 458 |
| $(b_{25})^{(4)}T_{24} - [(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}\left((G_{27})\right)]T_{25} = 0$ | 459 |
| $(b_{26})^{(4)}T_{25} - [(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}\left((G_{27})\right)]T_{26} = 0$ | 460 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29})\right]G_{28} = 0$ | 461 |
| $(a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29})\right]G_{29} = 0$ | 462 |
| $(a_{30})^{(5)}G_{29} - \left[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29})\right]G_{30} = 0$ | 463 |
| $(b_{28})^{(5)}T_{29} - [(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}(G_{31})]T_{28} = 0$ | 464 |
| $(b_{29})^{(5)}T_{28} - [(b_{29}^{'})^{(5)} - (b_{29}^{''})^{(5)}(G_{31})]T_{29} = 0$ | 465 |
| $(b_{30})^{(5)}T_{29} - [(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}(G_{31})]T_{30} = 0$ | 466 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{32})^{(6)}G_{33} - \left[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33})\right]G_{32} = 0$ | 467 |
| $(a_{33})^{(6)}G_{32} - \left[(a_{33}^{'})^{(6)} + (a_{33}^{''})^{(6)}(T_{33})\right]G_{33} = 0$ | 468 |
| $(a_{34})^{(6)}G_{33} - \left[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33})\right]G_{34} = 0$ | 469 |
| $(b_{32})^{(6)}T_{33} - [(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}(G_{35})]T_{32} = 0$ | 470 |
| $(b_{33})^{(6)}T_{32} - [(b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}(G_{35})]T_{33} = 0$ | 471 |
| $(b_{34})^{(6)}T_{33} - [(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}(G_{35})]T_{34} = 0$ | 472 |
| has a unique positive solution , which is an equilibrium solution | |

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| $(a_{36})^{(7)}G_{37} - \left[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37})\right]G_{36} = 0$ | 473 |
| $(a_{37})^{(7)}G_{36} - \left[(a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37})\right]G_{37} = 0$ | 474 |
| $(a_{38})^{(7)}G_{37} - \left[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37})\right]G_{38} = 0$ | 475 |
| $(b_{36})^{(7)}T_{37} - [(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}(G_{39})]T_{36} = 0$ | 476 |
| $(b_{37})^{(7)}T_{36} - [(b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}(G_{39})\,]T_{37} = 0$ | 477 |
| $(b_{38})^{(7)}T_{37} - [(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}(G_{39})\,]T_{38} = 0$ | 478 |
| $(a_{40})^{(8)}G_{41} - \left[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41})\right]G_{40} = 0$ | 479 |
| $(a_{41})^{(8)}G_{40} - \left[(a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41})\right]G_{41} = 0$ | 480 |
| $(a_{42})^{(8)}G_{41} - \left[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41})\right]G_{42} = 0$ | 481 |
| $(b_{40})^{(8)}T_{41} - [(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}(G_{43})\,]T_{40} = 0$ | 482 |
| $(b_{41})^{(8)}T_{40} - [(b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}(G_{43})\,]T_{41} = 0$ | 483 |
| $(b_{42})^{(8)}T_{41} - [(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}(G_{43})\,]T_{42} = 0$ | 484 |
| $(a_{44})^{(9)}G_{45} - \left[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45})\right]G_{44} = 0$ | 484 A |
| $(a_{45})^{(9)}G_{44} - \left[(a_{45}^{'})^{(9)} + (a_{45}^{''})^{(9)}(T_{45})\right]G_{45} = 0$ | |
| $(a_{46})^{(9)}G_{45} - \left[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45})\right]G_{46} = 0$ | |
| $(b_{44})^{(9)}T_{45} - [(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}(G_{47})\,]T_{44} = 0$ | |
| $(b_{45})^{(9)}T_{44} - [(b_{45}^{'})^{(9)} - (b_{45}^{''})^{(9)}(G_{47})\,]T_{45} = 0$ | |
| $(b_{46})^{(9)}T_{45} - [(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}(G_{47})\,]T_{46} = 0$ | |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{13}, G_{14}$ if<br><br>$F(T) = (a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a_{13}^{'})^{(1)}(a_{14}^{''})^{(1)}(T_{14}) + (a_{14}^{'})^{(1)}(a_{13}^{''})^{(1)}(T_{14}) +$ | 485 |

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| $(a_{13}^{''})^{(1)}(T_{14})(a_{14}^{''})^{(1)}(T_{14}) = 0$ | |
| **Proof:**<br><br>(c) Indeed the first two equations have a nontrivial solution $G_{16}, G_{17}$ if<br><br>$\mathrm{F}(T_{19}) = (a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a_{16}^{'})^{(2)}(a_{17}^{''})^{(2)}(T_{17}) + (a_{17}^{'})^{(2)}(a_{16}^{''})^{(2)}(T_{17}) + (a_{16}^{''})^{(2)}(T_{17})(a_{17}^{''})^{(2)}(T_{17}) = 0$ | 486 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{20}, G_{21}$ if<br><br>$F(T_{23}) = (a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a_{20}^{'})^{(3)}(a_{21}^{''})^{(3)}(T_{21}) + (a_{21}^{'})^{(3)}(a_{20}^{''})^{(3)}(T_{21}) + (a_{20}^{''})^{(3)}(T_{21})(a_{21}^{''})^{(3)}(T_{21}) = 0$ | 487 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{24}, G_{25}$ if<br><br>$F(T_{27}) = (a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a_{24}^{'})^{(4)}(a_{25}^{''})^{(4)}(T_{25}) + (a_{25}^{'})^{(4)}(a_{24}^{''})^{(4)}(T_{25}) + (a_{24}^{''})^{(4)}(T_{25})(a_{25}^{''})^{(4)}(T_{25}) = 0$ | 488 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{28}, G_{29}$ if<br><br>$F(T_{31}) = (a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a_{28}^{'})^{(5)}(a_{29}^{''})^{(5)}(T_{29}) + (a_{29}^{'})^{(5)}(a_{28}^{''})^{(5)}(T_{29}) + (a_{28}^{''})^{(5)}(T_{29})(a_{29}^{''})^{(5)}(T_{29}) = 0$ | 489 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{32}, G_{33}$ if<br><br>$F(T_{35}) = (a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32}^{'})^{(6)}(a_{33}^{''})^{(6)}(T_{33}) + (a_{33}^{'})^{(6)}(a_{32}^{''})^{(6)}(T_{33}) + (a_{32}^{''})^{(6)}(T_{33})(a_{33}^{''})^{(6)}(T_{33}) = 0$ | 490 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{36}, G_{37}$ if<br><br>$F(T_{39}) = (a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36}^{'})^{(7)}(a_{37}^{''})^{(7)}(T_{37}) + (a_{37}^{'})^{(7)}(a_{36}^{''})^{(7)}(T_{37}) + (a_{36}^{''})^{(7)}(T_{37})(a_{37}^{''})^{(7)}(T_{37}) = 0$ | 491 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{40}, G_{41}$ if<br><br>$F(T_{43}) = (a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40}^{'})^{(8)}(a_{41}^{''})^{(8)}(T_{41}) + (a_{41}^{'})^{(8)}(a_{40}^{''})^{(8)}(T_{41}) + (a_{40}^{''})^{(8)}(T_{41})(a_{41}^{''})^{(8)}(T_{41}) = 0$ | 492 |

**Proof:** 492 A

(a) Indeed the first two equations have a nontrivial solution $G_{44}, G_{45}$ if

$F(T_{47}) = (a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44}^{'})^{(9)}(a_{45}^{''})^{(9)}(T_{45}) + (a_{45}^{'})^{(9)}(a_{44}^{''})^{(9)}(T_{45}) + (a_{44}^{''})^{(9)}(T_{45})(a_{45}^{''})^{(9)}(T_{45}) = 0$

**Definition and uniqueness of** $T_{14}^{*}$ :- 493

After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{14})$ being increasing, it follows that there exists a unique $T_{14}^{*}$ for which $f(T_{14}^{*}) = 0$. With this value , we obtain from the three first equations

$G_{13} = \frac{(a_{13})^{(1)}G_{14}}{[(a_{13}^{'})^{(1)}+(a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15} = \frac{(a_{15})^{(1)}G_{14}}{[(a_{15}^{'})^{(1)}+(a_{15}^{''})^{(1)}(T_{14}^{*})]}$

**Definition and uniqueness of** $T_{17}^{*}$ :- 494

After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(2)}(T_{17})$ being increasing, it follows that there exists a unique $T_{17}^{*}$ for which $f(T_{17}^{*}) = 0$. With this value , we obtain from the three first equations

$G_{16} = \frac{(a_{16})^{(2)}G_{17}}{[(a_{16}^{'})^{(2)}+(a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18} = \frac{(a_{18})^{(2)}G_{17}}{[(a_{18}^{'})^{(2)}+(a_{18}^{''})^{(2)}(T_{17}^{*})]}$ 495

**Definition and uniqueness of** $T_{21}^{*}$ :- 496

After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{21})$ being increasing, it follows that there exists a unique $T_{21}^{*}$ for which $f(T_{21}^{*}) = 0$. With this value , we obtain from the three first equations

$G_{20} = \frac{(a_{20})^{(3)}G_{21}}{[(a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21}^{*})]}$ , $G_{22} = \frac{(a_{22})^{(3)}G_{21}}{[(a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21}^{*})]}$

**Definition and uniqueness of** $T_{25}^{*}$ :- 497

After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(4)}(T_{25})$ being increasing, it follows that there exists a unique $T_{25}^{*}$ for which $f(T_{25}^{*}) = 0$. With this value , we obtain from the three first equations

$G_{24} = \frac{(a_{24})^{(4)}G_{25}}{[(a_{24}^{'})^{(4)}+(a_{24}^{''})^{(4)}(T_{25}^{*})]}$ , $G_{26} = \frac{(a_{26})^{(4)}G_{25}}{[(a_{26}^{'})^{(4)}+(a_{26}^{''})^{(4)}(T_{25}^{*})]}$

**Definition and uniqueness of** $T_{29}^{*}$ :- 498

After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(5)}(T_{29})$ being increasing, it follows that there exists a unique $T_{29}^{*}$ for which $f(T_{29}^{*}) = 0$. With this value , we obtain from the three first equations

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| $G_{28} = \frac{(a_{28})^{(5)}G_{29}}{[(a'_{28})^{(5)}+(a''_{28})^{(5)}(T^*_{29})]}$ , $G_{30} = \frac{(a_{30})^{(5)}G_{29}}{[(a'_{30})^{(5)}+(a''_{30})^{(5)}(T^*_{29})]}$ | |
| **Definition and uniqueness of** $T^*_{33}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a''_i)^{(6)}(T_{33})$ being increasing, it follows that there exists a unique $T^*_{33}$ for which $f(T^*_{33}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{32} = \frac{(a_{32})^{(6)}G_{33}}{[(a'_{32})^{(6)}+(a''_{32})^{(6)}(T^*_{33})]}$ , $G_{34} = \frac{(a_{34})^{(6)}G_{33}}{[(a'_{34})^{(6)}+(a''_{34})^{(6)}(T^*_{33})]}$ | 499 |
| **Definition and uniqueness of** $T^*_{37}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a''_i)^{(7)}(T_{37})$ being increasing, it follows that there exists a unique $T^*_{37}$ for which $f(T^*_{37}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{36} = \frac{(a_{36})^{(7)}G_{37}}{[(a'_{36})^{(7)}+(a''_{36})^{(7)}(T^*_{37})]}$ , $G_{38} = \frac{(a_{38})^{(7)}G_{37}}{[(a'_{38})^{(7)}+(a''_{38})^{(7)}(T^*_{37})]}$ | 500 |
| **Definition and uniqueness of** $T^*_{41}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a''_i)^{(8)}(T_{41})$ being increasing, it follows that there exists a unique $T^*_{41}$ for which $f(T^*_{41}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{40} = \frac{(a_{40})^{(8)}G_{41}}{[(a'_{40})^{(8)}+(a''_{40})^{(8)}(T^*_{41})]}$ , $G_{42} = \frac{(a_{42})^{(8)}G_{41}}{[(a'_{42})^{(8)}+(a''_{42})^{(8)}(T^*_{41})]}$ | 501 |
| **Definition and uniqueness of** $T^*_{45}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a''_i)^{(9)}(T_{45})$ being increasing, it follows that there exists a unique $T^*_{45}$ for which $f(T^*_{45}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{44} = \frac{(a_{44})^{(9)}G_{45}}{[(a'_{44})^{(9)}+(a''_{44})^{(9)}(T^*_{45})]}$ , $G_{46} = \frac{(a_{46})^{(9)}G_{45}}{[(a'_{46})^{(9)}+(a''_{46})^{(9)}(T^*_{45})]}$ | 501 A |
| By the same argument, the equations admit solutions $G_{13}, G_{14}$ if<br>$\varphi(G) = (b'_{13})^{(1)}(b'_{14})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} -$<br><br>$\left[(b'_{13})^{(1)}(b''_{14})^{(1)}(G) + (b'_{14})^{(1)}(b''_{13})^{(1)}(G)\right]+(b''_{13})^{(1)}(G)(b''_{14})^{(1)}(G) = 0$<br><br>Where in $G(G_{13}, G_{14}, G_{15}), G_{13}, G_{15}$ must be replaced by their values from 96. It is easy to see that $\varphi$ is a decreasing function in $G_{14}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that | 502 |

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| there exists a unique $G_{14}^{*}$ such that $\varphi(G^{*})=0$ | |
| By the same argument, the equations admit solutions $G_{16}, G_{17}$ if<br>$\varphi(G_{19})=(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)}-(b_{16})^{(2)}(b_{17})^{(2)}-$<br>$\left[(b_{16}^{'})^{(2)}(b_{17}^{''})^{(2)}(G_{19})+(b_{17}^{'})^{(2)}(b_{16}^{''})^{(2)}(G_{19})\right]+(b_{16}^{''})^{(2)}(G_{19})(b_{17}^{''})^{(2)}(G_{19})=0$ | 503 |
| Where in $(G_{19})(G_{16},G_{17},G_{18}), G_{16}, G_{18}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{17}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{14}^{*}$ such that $\varphi((G_{19})^{*})=0$ | 504 |
| By the same argument, the equations admit solutions $G_{20}, G_{21}$ if<br>$\varphi(G_{23})=(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)}-(b_{20})^{(3)}(b_{21})^{(3)}-$<br>$\left[(b_{20}^{'})^{(3)}(b_{21}^{''})^{(3)}(G_{23})+(b_{21}^{'})^{(3)}(b_{20}^{''})^{(3)}(G_{23})\right]+(b_{20}^{''})^{(3)}(G_{23})(b_{21}^{''})^{(3)}(G_{23})=0$<br>Where in $G_{23}(G_{20},G_{21},G_{22}), G_{20}, G_{22}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{21}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{21}^{*}$ such that $\varphi((G_{23})^{*})=0$ | 505 |
| By the same argument, the equations admit solutions $G_{24}, G_{25}$ if<br>$\varphi(G_{27})=(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)}-(b_{24})^{(4)}(b_{25})^{(4)}-$<br>$\left[(b_{24}^{'})^{(4)}(b_{25}^{''})^{(4)}(G_{27})+(b_{25}^{'})^{(4)}(b_{24}^{''})^{(4)}(G_{27})\right]+(b_{24}^{''})^{(4)}(G_{27})(b_{25}^{''})^{(4)}(G_{27})=0$<br>Where in $(G_{27})(G_{24},G_{25},G_{26}), G_{24}, G_{26}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{25}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{25}^{*}$ such that $\varphi((G_{27})^{*})=0$ | 506 |
| By the same argument, the equations admit solutions $G_{28}, G_{29}$ if<br>$\varphi(G_{31})=(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)}-(b_{28})^{(5)}(b_{29})^{(5)}-$<br>$\left[(b_{28}^{'})^{(5)}(b_{29}^{''})^{(5)}(G_{31})+(b_{29}^{'})^{(5)}(b_{28}^{''})^{(5)}(G_{31})\right]+(b_{28}^{''})^{(5)}(G_{31})(b_{29}^{''})^{(5)}(G_{31})=0$<br>Where in $(G_{31})(G_{28},G_{29},G_{30}), G_{28}, G_{30}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{29}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{29}^{*}$ such that $\varphi((G_{31})^{*})=0$ | 507 |
| By the same argument, the equations admit solutions $G_{32}, G_{33}$ if<br>$\varphi(G_{35})=(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)}-(b_{32})^{(6)}(b_{33})^{(6)}-$<br>$\left[(b_{32}^{'})^{(6)}(b_{33}^{''})^{(6)}(G_{35})+(b_{33}^{'})^{(6)}(b_{32}^{''})^{(6)}(G_{35})\right]+(b_{32}^{''})^{(6)}(G_{35})(b_{33}^{''})^{(6)}(G_{35})=0$<br>Where in $(G_{35})(G_{32},G_{33},G_{34}), G_{32}, G_{34}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{33}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that | 508 |

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| there exists a unique $G_{33}^{*}$ such that $\varphi(G_{35}{}^{*}) = 0$ | |
| By the same argument, the equations admit solutions $G_{36}, G_{37}$ if<br>$\varphi(G_{39}) = (b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} -$<br>$\left[(b_{36}^{'})^{(7)}(b_{37}^{''})^{(7)}(G_{39}) + (b_{37}^{'})^{(7)}(b_{36}^{''})^{(7)}(G_{39})\right] + (b_{36}^{''})^{(7)}(G_{39})(b_{37}^{''})^{(7)}(G_{39}) = 0$<br><br>Where in $(G_{39})(G_{36}, G_{37}, G_{38}), G_{36}, G_{38}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{37}$ taking into account the hypothesis $\varphi(0) > 0\,, \varphi(\infty) < 0$ it follows that there exists a unique $G_{37}^{*}$ such that $\varphi(G_{39}{}^{*}) = 0$ | 509 |
| By the same argument, the equations admit solutions $G_{40}, G_{41}$ if<br>$\varphi(G_{43}) = (b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} -$<br>$\left[(b_{40}^{'})^{(8)}(b_{41}^{''})^{(8)}(G_{43}) + (b_{41}^{'})^{(8)}(b_{40}^{''})^{(8)}(G_{43})\right] + (b_{40}^{''})^{(8)}(G_{43})(b_{41}^{''})^{(8)}(G_{43}) = 0$<br><br>Where in $(G_{43})(G_{40}, G_{41}, G_{42}), G_{40}, G_{42}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{41}$ taking into account the hypothesis $\varphi(0) > 0\,, \varphi(\infty) < 0$ it follows that there exists a unique $G_{41}^{*}$ such that $\varphi(G_{43}{}^{*}) = 0$ | 510 |
| By the same argument, the equations 92,93 admit solutions $G_{44}, G_{45}$ if<br>$\varphi(G_{47}) = (b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} -$<br><br>$\left[(b_{44}^{'})^{(9)}(b_{45}^{''})^{(9)}(G_{47}) + (b_{45}^{'})^{(9)}(b_{44}^{''})^{(9)}(G_{47})\right] + (b_{44}^{''})^{(9)}(G_{47})(b_{45}^{''})^{(9)}(G_{47}) = 0$<br><br>Where in $(G_{47})(G_{44}, G_{45}, G_{46}), G_{44}, G_{46}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{45}$ taking into account the hypothesis $\varphi(0) > 0\,, \varphi(\infty) < 0$ it follows that there exists a unique $G_{45}^{*}$ such that $\varphi((G_{47})^{*}) = 0$ | |
| Finally we obtain the unique solution<br><br>$G_{14}^{*}$ given by $\varphi(G^{*}) = 0\,, T_{14}^{*}$ given by $f(T_{14}^{*}) = 0$ and<br><br>$G_{13}^{*} = \frac{(a_{13})^{(1)}G_{14}^{*}}{[(a_{13}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15}^{*} = \frac{(a_{15})^{(1)}G_{14}^{*}}{[(a_{15}^{'})^{(1)} + (a_{15}^{''})^{(1)}(T_{14}^{*})]}$<br><br>$T_{13}^{*} = \frac{(b_{13})^{(1)}T_{14}^{*}}{[(b_{13}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G^{*})]}$ , $T_{15}^{*} = \frac{(b_{15})^{(1)}T_{14}^{*}}{[(b_{15}^{'})^{(1)} - (b_{15}^{''})^{(1)}(G^{*})]}$<br><br>Obviously, these values represent an equilibrium solution | 511 |
| Finally we obtain the unique solution | |
| $G_{17}^{*}$ given by $\varphi((G_{19})^{*}) = 0\,, T_{17}^{*}$ given by $f(T_{17}^{*}) = 0$ and | 512 |
| $G_{16}^{*} = \frac{(a_{16})^{(2)}G_{17}^{*}}{[(a_{16}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18}^{*} = \frac{(a_{18})^{(2)}G_{17}^{*}}{[(a_{18}^{'})^{(2)} + (a_{18}^{''})^{(2)}(T_{17}^{*})]}$ | 513 |
| $T_{16}^{*} = \frac{(b_{16})^{(2)}T_{17}^{*}}{[(b_{16}^{'})^{(2)} - (b_{16}^{''})^{(2)}((G_{19})^{*})]}$ , $T_{18}^{*} = \frac{(b_{18})^{(2)}T_{17}^{*}}{[(b_{18}^{'})^{(2)} - (b_{18}^{''})^{(2)}((G_{19})^{*})]}$ | 514 |

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| Obviously, these values represent an equilibrium solution | |
| Finally we obtain the unique solution<br><br>$G_{21}^*$ given by $\varphi((G_{23})^*) = 0$ , $T_{21}^*$ given by $f(T_{21}^*) = 0$ and<br><br>$G_{20}^* = \frac{(a_{20})^{(3)} G_{21}^*}{[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21}^*)]}$ , $G_{22}^* = \frac{(a_{22})^{(3)} G_{21}^*}{[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21}^*)]}$<br><br>$T_{20}^* = \frac{(b_{20})^{(3)} T_{21}^*}{[(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23}{}^*)]}$ , $T_{22}^* = \frac{(b_{22})^{(3)} T_{21}^*}{[(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23}{}^*)]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 515 |
| Finally we obtain the unique solution<br><br>$G_{25}^*$ given by $\varphi(G_{27}) = 0$ , $T_{25}^*$ given by $f(T_{25}^*) = 0$ and<br><br>$G_{24}^* = \frac{(a_{24})^{(4)} G_{25}^*}{[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25}^*)]}$ , $G_{26}^* = \frac{(a_{26})^{(4)} G_{25}^*}{[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25}^*)]}$ | 516 |
| $T_{24}^* = \frac{(b_{24})^{(4)} T_{25}^*}{[(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}((G_{27})^*)]}$ , $T_{26}^* = \frac{(b_{26})^{(4)} T_{25}^*}{[(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}((G_{27})^*)]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 517 |
| Finally we obtain the unique solution<br><br>$G_{29}^*$ given by $\varphi((G_{31})^*) = 0$ , $T_{29}^*$ given by $f(T_{29}^*) = 0$ and<br><br>$G_{28}^* = \frac{(a_{28})^{(5)} G_{29}^*}{[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}^*)]}$ , $G_{30}^* = \frac{(a_{30})^{(5)} G_{29}^*}{[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29}^*)]}$ | 518 |
| $T_{28}^* = \frac{(b_{28})^{(5)} T_{29}^*}{[(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}((G_{31})^*)]}$ , $T_{30}^* = \frac{(b_{30})^{(5)} T_{29}^*}{[(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}((G_{31})^*)]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 519 |
| Finally we obtain the unique solution<br><br>$G_{33}^*$ given by $\varphi((G_{35})^*) = 0$ , $T_{33}^*$ given by $f(T_{33}^*) = 0$ and<br><br>$G_{32}^* = \frac{(a_{32})^{(6)} G_{33}^*}{[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33}^*)]}$ , $G_{34}^* = \frac{(a_{34})^{(6)} G_{33}^*}{[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33}^*)]}$ | 520 |
| $T_{32}^* = \frac{(b_{32})^{(6)} T_{33}^*}{[(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}((G_{35})^*)]}$ , $T_{34}^* = \frac{(b_{34})^{(6)} T_{33}^*}{[(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}((G_{35})^*)]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 521 |
| Finally we obtain the unique solution<br><br>$G_{37}^*$ given by $\varphi((G_{39})^*) = 0$ , $T_{37}^*$ given by $f(T_{37}^*) = 0$ and | 522 |

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| $G_{36}^{*} = \frac{(a_{36})^{(7)} G_{37}^{*}}{[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)} (T_{37}^{*})]}$ , $G_{38}^{*} = \frac{(a_{38})^{(7)} G_{37}^{*}}{[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)} (T_{37}^{*})]}$<br><br>$T_{36}^{*} = \frac{(b_{36})^{(7)} T_{37}^{*}}{[(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)} ((G_{39})^{*})]}$ , $T_{38}^{*} = \frac{(b_{38})^{(7)} T_{37}^{*}}{[(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)} ((G_{39})^{*})]}$ | |
| Finally we obtain the unique solution<br><br>$G_{41}^{*}$ given by $\varphi((G_{43})^{*}) = 0$ , $T_{41}^{*}$ given by $f(T_{41}^{*}) = 0$ and<br><br>$G_{40}^{*} = \frac{(a_{40})^{(8)} G_{41}^{*}}{[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)} (T_{41}^{*})]}$ , $G_{42}^{*} = \frac{(a_{42})^{(8)} G_{41}^{*}}{[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)} (T_{41}^{*})]}$<br><br>$T_{40}^{*} = \frac{(b_{40})^{(8)} T_{41}^{*}}{[(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)} ((G_{43})^{*})]}$ , $T_{42}^{*} = \frac{(b_{42})^{(8)} T_{41}^{*}}{[(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)} ((G_{43})^{*})]}$ | 523 |
| Finally we obtain the unique solution of 89 to 99<br><br>$G_{45}^{*}$ given by $\varphi((G_{47})^{*}) = 0$ , $T_{45}^{*}$ given by $f(T_{45}^{*}) = 0$ and<br><br>$G_{44}^{*} = \frac{(a_{44})^{(9)} G_{45}^{*}}{[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)} (T_{45}^{*})]}$ , $G_{46}^{*} = \frac{(a_{46})^{(9)} G_{45}^{*}}{[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)} (T_{45}^{*})]}$<br><br>$T_{44}^{*} = \frac{(b_{44})^{(9)} T_{45}^{*}}{[(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)} ((G_{47})^{*})]}$ , $T_{46}^{*} = \frac{(b_{46})^{(9)} T_{45}^{*}}{[(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)} ((G_{47})^{*})]}$ | 523 A |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(1)}$ *and* $(b_i^{''})^{(1)}$ Belong to $C^{(1)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>**Proof:**_Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^{*} + \mathbb{G}_i$ , $T_i = T_i^{*} + \mathbb{T}_i$<br><br>$\frac{\partial (a_{14}^{''})^{(1)}}{\partial T_{14}} (T_{14}^{*}) = (q_{14})^{(1)}$ , $\frac{\partial (b_i^{''})^{(1)}}{\partial G_j} (G^{*}) = s_{ij}$ | 524 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{13}}{dt} = -\left((a_{13}^{'})^{(1)} + (p_{13})^{(1)}\right)\mathbb{G}_{13} + (a_{13})^{(1)}\mathbb{G}_{14} - (q_{13})^{(1)} G_{13}^{*} \mathbb{T}_{14}$ | 525 |

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| $\frac{d\mathbb{G}_{14}}{dt} = -\left((a_{14}^{'})^{(1)} + (p_{14})^{(1)}\right)\mathbb{G}_{14} + (a_{14})^{(1)}\mathbb{G}_{13} - (q_{14})^{(1)}G_{14}^{*}\mathbb{T}_{14}$ | 526 |
| $\frac{d\mathbb{G}_{15}}{dt} = -\left((a_{15}^{'})^{(1)} + (p_{15})^{(1)}\right)\mathbb{G}_{15} + (a_{15})^{(1)}\mathbb{G}_{14} - (q_{15})^{(1)}G_{15}^{*}\mathbb{T}_{14}$ | 527 |
| $\frac{d\mathbb{T}_{13}}{dt} = -\left((b_{13}^{'})^{(1)} - (r_{13})^{(1)}\right)\mathbb{T}_{13} + (b_{13})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(13)(j)}T_{13}^{*}\mathbb{G}_{j}\right)$ | 528 |
| $\frac{d\mathbb{T}_{14}}{dt} = -\left((b_{14}^{'})^{(1)} - (r_{14})^{(1)}\right)\mathbb{T}_{14} + (b_{14})^{(1)}\mathbb{T}_{13} + \sum_{j=13}^{15}\left(s_{(14)(j)}T_{14}^{*}\mathbb{G}_{j}\right)$ | 529 |
| $\frac{d\mathbb{T}_{15}}{dt} = -\left((b_{15}^{'})^{(1)} - (r_{15})^{(1)}\right)\mathbb{T}_{15} + (b_{15})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(15)(j)}T_{15}^{*}\mathbb{G}_{j}\right)$ | 530 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}$ Belong to $C^{(2)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable | 531 |
| Proof: Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- | |
| $G_i = G_i^{*} + \mathbb{G}_i \quad , T_i = T_i^{*} + \mathbb{T}_i$ | 532 |
| $\frac{\partial(a_{17}^{''})^{(2)}}{\partial T_{17}}(T_{17}^{*}) = (q_{17})^{(2)}$ , $\frac{\partial(b_i^{''})^{(2)}}{\partial G_j}((G_{19})^{*}) = s_{ij}$ | 533 |
| taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{16}}{dt} = -\left((a_{16}^{'})^{(2)} + (p_{16})^{(2)}\right)\mathbb{G}_{16} + (a_{16})^{(2)}\mathbb{G}_{17} - (q_{16})^{(2)}G_{16}^{*}\mathbb{T}_{17}$ | 534 |
| $\frac{d\mathbb{G}_{17}}{dt} = -\left((a_{17}^{'})^{(2)} + (p_{17})^{(2)}\right)\mathbb{G}_{17} + (a_{17})^{(2)}\mathbb{G}_{16} - (q_{17})^{(2)}G_{17}^{*}\mathbb{T}_{17}$ | 535 |
| $\frac{d\mathbb{G}_{18}}{dt} = -\left((a_{18}^{'})^{(2)} + (p_{18})^{(2)}\right)\mathbb{G}_{18} + (a_{18})^{(2)}\mathbb{G}_{17} - (q_{18})^{(2)}G_{18}^{*}\mathbb{T}_{17}$ | 536 |
| $\frac{d\mathbb{T}_{16}}{dt} = -\left((b_{16}^{'})^{(2)} - (r_{16})^{(2)}\right)\mathbb{T}_{16} + (b_{16})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(16)(j)}T_{16}^{*}\mathbb{G}_{j}\right)$ | 537 |
| $\frac{d\mathbb{T}_{17}}{dt} = -\left((b_{17}^{'})^{(2)} - (r_{17})^{(2)}\right)\mathbb{T}_{17} + (b_{17})^{(2)}\mathbb{T}_{16} + \sum_{j=16}^{18}\left(s_{(17)(j)}T_{17}^{*}\mathbb{G}_{j}\right)$ | 538 |
| $\frac{d\mathbb{T}_{18}}{dt} = -\left((b_{18}^{'})^{(2)} - (r_{18})^{(2)}\right)\mathbb{T}_{18} + (b_{18})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(18)(j)}T_{18}^{*}\mathbb{G}_{j}\right)$ | 539 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(3)}$ *and* $(b_i^{''})^{(3)}$ Belong to $C^{(3)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- | 540 |

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| $G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br>$\frac{\partial (a_{21}^{''})^{(3)}}{\partial T_{21}}(T_{21}^*) = (q_{21})^{(3)} \ , \ \frac{\partial (b_i^{''})^{(3)}}{\partial G_j}((G_{23})^*) = s_{ij}$ | |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{20}}{dt} = -\left((a_{20}^{'})^{(3)} + (p_{20})^{(3)}\right)\mathbb{G}_{20} + (a_{20})^{(3)}\mathbb{G}_{21} - (q_{20})^{(3)}G_{20}^*\mathbb{T}_{21}$ | 541 |
| $\frac{d\mathbb{G}_{21}}{dt} = -\left((a_{21}^{'})^{(3)} + (p_{21})^{(3)}\right)\mathbb{G}_{21} + (a_{21})^{(3)}\mathbb{G}_{20} - (q_{21})^{(3)}G_{21}^*\mathbb{T}_{21}$ | 542 |
| $\frac{d\mathbb{G}_{22}}{dt} = -\left((a_{22}^{'})^{(3)} + (p_{22})^{(3)}\right)\mathbb{G}_{22} + (a_{22})^{(3)}\mathbb{G}_{21} - (q_{22})^{(3)}G_{22}^*\mathbb{T}_{21}$ | 543 |
| $\frac{d\mathbb{T}_{20}}{dt} = -\left((b_{20}^{'})^{(3)} - (r_{20})^{(3)}\right)\mathbb{T}_{20} + (b_{20})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(20)(j)}T_{20}^*\mathbb{G}_j\right)$ | 544 |
| $\frac{d\mathbb{T}_{21}}{dt} = -\left((b_{21}^{'})^{(3)} - (r_{21})^{(3)}\right)\mathbb{T}_{21} + (b_{21})^{(3)}\mathbb{T}_{20} + \sum_{j=20}^{22}\left(s_{(21)(j)}T_{21}^*\mathbb{G}_j\right)$ | 545 |
| $\frac{d\mathbb{T}_{22}}{dt} = -\left((b_{22}^{'})^{(3)} - (r_{22})^{(3)}\right)\mathbb{T}_{22} + (b_{22})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(22)(j)}T_{22}^*\mathbb{G}_j\right)$ | 546 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(4)}$ *and* $(b_i^{''})^{(4)}$ Belong to $C^{(4)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 547 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br>$\frac{\partial (a_{25}^{''})^{(4)}}{\partial T_{25}}(T_{25}^*) = (q_{25})^{(4)} \ , \ \frac{\partial (b_i^{''})^{(4)}}{\partial G_j}((G_{27})^*) = s_{ij}$ | 548 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{24}}{dt} = -\left((a_{24}^{'})^{(4)} + (p_{24})^{(4)}\right)\mathbb{G}_{24} + (a_{24})^{(4)}\mathbb{G}_{25} - (q_{24})^{(4)}G_{24}^*\mathbb{T}_{25}$ | 549 |
| $\frac{d\mathbb{G}_{25}}{dt} = -\left((a_{25}^{'})^{(4)} + (p_{25})^{(4)}\right)\mathbb{G}_{25} + (a_{25})^{(4)}\mathbb{G}_{24} - (q_{25})^{(4)}G_{25}^*\mathbb{T}_{25}$ | 550 |
| $\frac{d\mathbb{G}_{26}}{dt} = -\left((a_{26}^{'})^{(4)} + (p_{26})^{(4)}\right)\mathbb{G}_{26} + (a_{26})^{(4)}\mathbb{G}_{25} - (q_{26})^{(4)}G_{26}^*\mathbb{T}_{25}$ | 551 |
| $\frac{d\mathbb{T}_{24}}{dt} = -\left((b_{24}^{'})^{(4)} - (r_{24})^{(4)}\right)\mathbb{T}_{24} + (b_{24})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(24)(j)}T_{24}^*\mathbb{G}_j\right)$ | 552 |
| $\frac{d\mathbb{T}_{25}}{dt} = -\left((b_{25}^{'})^{(4)} - (r_{25})^{(4)}\right)\mathbb{T}_{25} + (b_{25})^{(4)}\mathbb{T}_{24} + \sum_{j=24}^{26}\left(s_{(25)(j)}T_{25}^*\mathbb{G}_j\right)$ | 553 |
| $\frac{d\mathbb{T}_{26}}{dt} = -\left((b_{26}^{'})^{(4)} - (r_{26})^{(4)}\right)\mathbb{T}_{26} + (b_{26})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(26)(j)}T_{26}^*\mathbb{G}_j\right)$ | 554 |

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| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 5:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(5)}$ *and* $(b_i^{''})^{(5)}$ Belong to $C^{(5)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 555 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial (a_{29}^{''})^{(5)}}{\partial T_{29}}(T_{29}^*) = (q_{29})^{(5)}$ , $\frac{\partial (b_i^{''})^{(5)}}{\partial G_j}((G_{31})^*) = s_{ij}$ | 556 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{28}}{dt} = -\left((a_{28}^{'})^{(5)} + (p_{28})^{(5)}\right)\mathbb{G}_{28} + (a_{28})^{(5)}\mathbb{G}_{29} - (q_{28})^{(5)}G_{28}^*\mathbb{T}_{29}$ | 557 |
| $\frac{d\mathbb{G}_{29}}{dt} = -\left((a_{29}^{'})^{(5)} + (p_{29})^{(5)}\right)\mathbb{G}_{29} + (a_{29})^{(5)}\mathbb{G}_{28} - (q_{29})^{(5)}G_{29}^*\mathbb{T}_{29}$ | 558 |
| $\frac{d\mathbb{G}_{30}}{dt} = -\left((a_{30}^{'})^{(5)} + (p_{30})^{(5)}\right)\mathbb{G}_{30} + (a_{30})^{(5)}\mathbb{G}_{29} - (q_{30})^{(5)}G_{30}^*\mathbb{T}_{29}$ | 559 |
| $\frac{d\mathbb{T}_{28}}{dt} = -\left((b_{28}^{'})^{(5)} - (r_{28})^{(5)}\right)\mathbb{T}_{28} + (b_{28})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(28)(j)}T_{28}^*\mathbb{G}_j\right)$ | 560 |
| $\frac{d\mathbb{T}_{29}}{dt} = -\left((b_{29}^{'})^{(5)} - (r_{29})^{(5)}\right)\mathbb{T}_{29} + (b_{29})^{(5)}\mathbb{T}_{28} + \sum_{j=28}^{30}\left(s_{(29)(j)}T_{29}^*\mathbb{G}_j\right)$ | 561 |
| $\frac{d\mathbb{T}_{30}}{dt} = -\left((b_{30}^{'})^{(5)} - (r_{30})^{(5)}\right)\mathbb{T}_{30} + (b_{30})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(30)(j)}T_{30}^*\mathbb{G}_j\right)$ | 562 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 6:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(6)}$ *and* $(b_i^{''})^{(6)}$ Belong to $C^{(6)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 563 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial (a_{33}^{''})^{(6)}}{\partial T_{33}}(T_{33}^*) = (q_{33})^{(6)}$ , $\frac{\partial (b_i^{''})^{(6)}}{\partial G_j}((G_{35})^*) = s_{ij}$ | 564 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{32}}{dt} = -\left((a_{32}^{'})^{(6)} + (p_{32})^{(6)}\right)\mathbb{G}_{32} + (a_{32})^{(6)}\mathbb{G}_{33} - (q_{32})^{(6)}G_{32}^*\mathbb{T}_{33}$ | 565 |
| $\frac{d\mathbb{G}_{33}}{dt} = -\left((a_{33}^{'})^{(6)} + (p_{33})^{(6)}\right)\mathbb{G}_{33} + (a_{33})^{(6)}\mathbb{G}_{32} - (q_{33})^{(6)}G_{33}^*\mathbb{T}_{33}$ | 566 |

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| $\frac{d\mathbb{G}_{34}}{dt} = -\left((a'_{34})^{(6)} + (p_{34})^{(6)}\right)\mathbb{G}_{34} + (a_{34})^{(6)}\mathbb{G}_{33} - (q_{34})^{(6)}G^*_{34}\mathbb{T}_{33}$ | 567 |
| $\frac{d\mathbb{T}_{32}}{dt} = -\left((b'_{32})^{(6)} - (r_{32})^{(6)}\right)\mathbb{T}_{32} + (b_{32})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(32)(j)}T^*_{32}\mathbb{G}_j\right)$ | 568 |
| $\frac{d\mathbb{T}_{33}}{dt} = -\left((b'_{33})^{(6)} - (r_{33})^{(6)}\right)\mathbb{T}_{33} + (b_{33})^{(6)}\mathbb{T}_{32} + \sum_{j=32}^{34}\left(s_{(33)(j)}T^*_{33}\mathbb{G}_j\right)$ | 569 |
| $\frac{d\mathbb{T}_{34}}{dt} = -\left((b'_{34})^{(6)} - (r_{34})^{(6)}\right)\mathbb{T}_{34} + (b_{34})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(34)(j)}T^*_{34}\mathbb{G}_j\right)$ | 570 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 7:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(7)}$ *and* $(b_i^{''})^{(7)}$ Belong to $C^{(7)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 571 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial(a''_{37})^{(7)}}{\partial T_{37}}(T^*_{37}) = (q_{37})^{(7)}$ , $\frac{\partial(b_i^{''})^{(7)}}{\partial G_j}\left((G_{39})^{**}\right) = s_{ij}$ | 572 |
| Then taking into account equations and neglecting the terms of power 2, we obtain from | |
| $\frac{d\mathbb{G}_{36}}{dt} = -\left((a'_{36})^{(7)} + (p_{36})^{(7)}\right)\mathbb{G}_{36} + (a_{36})^{(7)}\mathbb{G}_{37} - (q_{36})^{(7)}G^*_{36}\mathbb{T}_{37}$ | 573 |
| $\frac{d\mathbb{G}_{37}}{dt} = -\left((a'_{37})^{(7)} + (p_{37})^{(7)}\right)\mathbb{G}_{37} + (a_{37})^{(7)}\mathbb{G}_{36} - (q_{37})^{(7)}G^*_{37}\mathbb{T}_{37}$ | 574 |
| $\frac{d\mathbb{G}_{38}}{dt} = -\left((a'_{38})^{(7)} + (p_{38})^{(7)}\right)\mathbb{G}_{38} + (a_{38})^{(7)}\mathbb{G}_{37} - (q_{38})^{(7)}G^*_{38}\mathbb{T}_{37}$ | 575 |
| $\frac{d\mathbb{T}_{36}}{dt} = -\left((b'_{36})^{(7)} - (r_{36})^{(7)}\right)\mathbb{T}_{36} + (b_{36})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(36)(j)}T^*_{36}\mathbb{G}_j\right)$ | 576 |
| $\frac{d\mathbb{T}_{37}}{dt} = -\left((b'_{37})^{(7)} - (r_{37})^{(7)}\right)\mathbb{T}_{37} + (b_{37})^{(7)}\mathbb{T}_{36} + \sum_{j=36}^{38}\left(s_{(37)(j)}T^*_{37}\mathbb{G}_j\right)$ | 578 |
| $\frac{d\mathbb{T}_{38}}{dt} = -\left((b'_{38})^{(7)} - (r_{38})^{(7)}\right)\mathbb{T}_{38} + (b_{38})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(38)(j)}T^*_{38}\mathbb{G}_j\right)$ | 579 |
| Obviously, these values represent an equilibrium solution<br><br>**ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 8:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(8)}$ *and* $(b_i^{''})^{(8)}$ Belong to $C^{(8)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable. | |

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| Proof: Denote | |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- $G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$ $\frac{\partial (a_{41}^{''})^{(8)}}{\partial T_{41}}(T_{41}^*) = (q_{41})^{(8)} \ , \frac{\partial (b_i^{''})^{(8)}}{\partial G_j}((G_{43})^*) = s_{ij}$ | 580 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{40}}{dt} = -\left((a_{40}^{'})^{(8)} + (p_{40})^{(8)}\right)\mathbb{G}_{40} + (a_{40})^{(8)}\mathbb{G}_{41} - (q_{40})^{(8)}G_{40}^*\mathbb{T}_{41}$ | 581 |
| $\frac{d\mathbb{G}_{41}}{dt} = -\left((a_{41}^{'})^{(8)} + (p_{41})^{(8)}\right)\mathbb{G}_{41} + (a_{41})^{(8)}\mathbb{G}_{40} - (q_{41})^{(8)}G_{41}^*\mathbb{T}_{41}$ | 582 |
| $\frac{d\mathbb{G}_{42}}{dt} = -\left((a_{42}^{'})^{(8)} + (p_{42})^{(8)}\right)\mathbb{G}_{42} + (a_{42})^{(8)}\mathbb{G}_{41} - (q_{42})^{(8)}G_{42}^*\mathbb{T}_{41}$ | 583 |
| $\frac{d\mathbb{T}_{40}}{dt} = -\left((b_{40}^{'})^{(8)} - (r_{40})^{(8)}\right)\mathbb{T}_{40} + (b_{40})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(40)(j)}T_{40}^*\mathbb{G}_j\right)$ | 584 |
| $\frac{d\mathbb{T}_{41}}{dt} = -\left((b_{41}^{'})^{(8)} - (r_{41})^{(8)}\right)\mathbb{T}_{41} + (b_{41})^{(8)}\mathbb{T}_{40} + \sum_{j=40}^{42}\left(s_{(41)(j)}T_{41}^*\mathbb{G}_j\right)$ | 585 |
| $\frac{d\mathbb{T}_{42}}{dt} = -\left((b_{42}^{'})^{(8)} - (r_{42})^{(8)}\right)\mathbb{T}_{42} + (b_{42})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(42)(j)}T_{42}^*\mathbb{G}_j\right)$ | 586 |
| **ASYMPTOTIC STABILITY ANALYSIS** **Theorem 9:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(9)}$ *and* $(b_i^{''})^{(9)}$ Belong to $C^{(9)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable. **Proof: Denote** | 586 A |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- $G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$ $\frac{\partial (a_{45}^{''})^{(9)}}{\partial T_{45}}(T_{45}^*) = (q_{45})^{(9)} \ , \frac{\partial (b_i^{''})^{(9)}}{\partial G_j}((G_{47})^*) = s_{ij}$ | |
| Then taking into account equations 89 to 99 and neglecting the terms of power 2, we obtain from 99 to 44 | |
| $\frac{d\mathbb{G}_{44}}{dt} = -\left((a_{44}^{'})^{(9)} + (p_{44})^{(9)}\right)\mathbb{G}_{44} + (a_{44})^{(9)}\mathbb{G}_{45} - (q_{44})^{(9)}G_{44}^*\mathbb{T}_{45}$ | 586 B |
| $\frac{d\mathbb{G}_{45}}{dt} = -\left((a_{45}^{'})^{(9)} + (p_{45})^{(9)}\right)\mathbb{G}_{45} + (a_{45})^{(9)}\mathbb{G}_{44} - (q_{45})^{(9)}G_{45}^*\mathbb{T}_{45}$ | 586 C |
| $\frac{d\mathbb{G}_{46}}{dt} = -\left((a_{46}^{'})^{(9)} + (p_{46})^{(9)}\right)\mathbb{G}_{46} + (a_{46})^{(9)}\mathbb{G}_{45} - (q_{46})^{(9)}G_{46}^*\mathbb{T}_{45}$ | 586 D |

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| $\frac{d\mathbb{T}_{44}}{dt} = -\left((b'_{44})^{(9)} - (r_{44})^{(9)}\right)\mathbb{T}_{44} + (b_{44})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(44)(j)}T^*_{44}\mathbb{G}_j\right)$ | 586 E |
| $\frac{d\mathbb{T}_{45}}{dt} = -\left((b'_{45})^{(9)} - (r_{45})^{(9)}\right)\mathbb{T}_{45} + (b_{45})^{(9)}\mathbb{T}_{44} + \sum_{j=44}^{46}\left(s_{(45)(j)}T^*_{45}\mathbb{G}_j\right)$ | 586 F |
| $\frac{d\mathbb{T}_{46}}{dt} = -\left((b'_{46})^{(9)} - (r_{46})^{(9)}\right)\mathbb{T}_{46} + (b_{46})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(46)(j)}T^*_{46}\mathbb{G}_j\right)$ | 586 G |
| **The characteristic equation of this system is** | 587 |
| $\left((\lambda)^{(1)} + (b'_{15})^{(1)} - (r_{15})^{(1)}\right)\{\left((\lambda)^{(1)} + (a'_{15})^{(1)} + (p_{15})^{(1)}\right)$ $\left[\left(\left((\lambda)^{(1)} + (a'_{13})^{(1)} + (p_{13})^{(1)}\right)(q_{14})^{(1)}G^*_{14} + (a_{14})^{(1)}(q_{13})^{(1)}G^*_{13}\right)\right]$ $\left(\left((\lambda)^{(1)} + (b'_{13})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(14)}T^*_{14} + (b_{14})^{(1)}s_{(13),(14)}T^*_{14}\right)$ $+\left(\left((\lambda)^{(1)} + (a'_{14})^{(1)} + (p_{14})^{(1)}\right)(q_{13})^{(1)}G^*_{13} + (a_{13})^{(1)}(q_{14})^{(1)}G^*_{14}\right)$ $\left(\left((\lambda)^{(1)} + (b'_{13})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(13)}T^*_{14} + (b_{14})^{(1)}s_{(13),(13)}T^*_{13}\right)$ $\left(\left((\lambda)^{(1)}\right)^2 + \left((a'_{13})^{(1)} + (a'_{14})^{(1)} + (p_{13})^{(1)} + (p_{14})^{(1)}\right)(\lambda)^{(1)}\right)$ $\left(\left((\lambda)^{(1)}\right)^2 + \left((b'_{13})^{(1)} + (b'_{14})^{(1)} - (r_{13})^{(1)} + (r_{14})^{(1)}\right)(\lambda)^{(1)}\right)$ $+\left(\left((\lambda)^{(1)}\right)^2 + \left((a'_{13})^{(1)} + (a'_{14})^{(1)} + (p_{13})^{(1)} + (p_{14})^{(1)}\right)(\lambda)^{(1)}\right)(q_{15})^{(1)}G_{15}$ $+\left((\lambda)^{(1)} + (a'_{13})^{(1)} + (p_{13})^{(1)}\right)\left((a_{15})^{(1)}(q_{14})^{(1)}G^*_{14} + (a_{14})^{(1)}(a_{15})^{(1)}(q_{13})^{(1)}G^*_{13}\right)$ $\left(\left((\lambda)^{(1)} + (b'_{13})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(15)}T^*_{14} + (b_{14})^{(1)}s_{(13),(15)}T^*_{13}\right)\} = 0$ $+$ $\left((\lambda)^{(2)} + (b'_{18})^{(2)} - (r_{18})^{(2)}\right)\{\left((\lambda)^{(2)} + (a'_{18})^{(2)} + (p_{18})^{(2)}\right)$ $\left[\left(\left((\lambda)^{(2)} + (a'_{16})^{(2)} + (p_{16})^{(2)}\right)(q_{17})^{(2)}\mathrm{G}^*_{17} + (a_{17})^{(2)}(q_{16})^{(2)}\mathrm{G}^*_{16}\right)\right]$ $\left(\left((\lambda)^{(2)} + (b'_{16})^{(2)} - (r_{16})^{(2)}\right)s_{(17),(17)}\mathrm{T}^*_{17} + (b_{17})^{(2)}s_{(16),(17)}\mathrm{T}^*_{17}\right)$ $+\left(\left((\lambda)^{(2)} + (a'_{17})^{(2)} + (p_{17})^{(2)}\right)(q_{16})^{(2)}\mathrm{G}^*_{16} + (a_{16})^{(2)}(q_{17})^{(2)}\mathrm{G}^*_{17}\right)$ $\left(\left((\lambda)^{(2)} + (b'_{16})^{(2)} - (r_{16})^{(2)}\right)s_{(17),(16)}\mathrm{T}^*_{17} + (b_{17})^{(2)}s_{(16),(16)}\mathrm{T}^*_{16}\right)$ $\left(\left((\lambda)^{(2)}\right)^2 + \left((a'_{16})^{(2)} + (a'_{17})^{(2)} + (p_{16})^{(2)} + (p_{17})^{(2)}\right)(\lambda)^{(2)}\right)$ $\left(\left((\lambda)^{(2)}\right)^2 + \left((b'_{16})^{(2)} + (b'_{17})^{(2)} - (r_{16})^{(2)} + (r_{17})^{(2)}\right)(\lambda)^{(2)}\right)$ | |

$$+\left(\left((\lambda)^{(2)}\right)^2+\left((a_{16}')^{(2)}+(a_{17}')^{(2)}+(p_{16})^{(2)}+(p_{17})^{(2)}\right)(\lambda)^{(2)}\right)(q_{18})^{(2)}\mathrm{G}_{18}$$

$$+\left((\lambda)^{(2)}+(a_{16}')^{(2)}+(p_{16})^{(2)}\right)\left((a_{18})^{(2)}(q_{17})^{(2)}\mathrm{G}_{17}^{*}+(a_{17})^{(2)}(a_{18})^{(2)}(q_{16})^{(2)}\mathrm{G}_{16}^{*}\right)$$

$$\left(\left((\lambda)^{(2)}+(b_{16}')^{(2)}-(r_{16})^{(2)}\right)s_{(17),(18)}\mathrm{T}_{17}^{*}+(b_{17})^{(2)}s_{(16),(18)}\mathrm{T}_{16}^{*}\right)\}=0$$

$$+$$

$$\left((\lambda)^{(3)}+(b_{22}')^{(3)}-(r_{22})^{(3)}\right)\{\left((\lambda)^{(3)}+(a_{22}')^{(3)}+(p_{22})^{(3)}\right)$$

$$\left[\left(\left((\lambda)^{(3)}+(a_{20}')^{(3)}+(p_{20})^{(3)}\right)(q_{21})^{(3)}G_{21}^{*}+(a_{21})^{(3)}(q_{20})^{(3)}G_{20}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(21)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(21)}T_{21}^{*}\right)$$

$$+\left(\left((\lambda)^{(3)}+(a_{21}')^{(3)}+(p_{21})^{(3)}\right)(q_{20})^{(3)}G_{20}^{*}+(a_{20})^{(3)}(q_{21})^{(1)}G_{21}^{*}\right)$$

$$\left(\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(20)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(20)}T_{20}^{*}\right)$$

$$\left(\left((\lambda)^{(3)}\right)^2+\left((a_{20}')^{(3)}+(a_{21}')^{(3)}+(p_{20})^{(3)}+(p_{21})^{(3)}\right)(\lambda)^{(3)}\right)$$

$$\left(\left((\lambda)^{(3)}\right)^2+\left((b_{20}')^{(3)}+(b_{21}')^{(3)}-(r_{20})^{(3)}+(r_{21})^{(3)}\right)(\lambda)^{(3)}\right)$$

$$+\left(\left((\lambda)^{(3)}\right)^2+\left((a_{20}')^{(3)}+(a_{21}')^{(3)}+(p_{20})^{(3)}+(p_{21})^{(3)}\right)(\lambda)^{(3)}\right)(q_{22})^{(3)}G_{22}$$

$$+\left((\lambda)^{(3)}+(a_{20}')^{(3)}+(p_{20})^{(3)}\right)\left((a_{22})^{(3)}(q_{21})^{(3)}G_{21}^{*}+(a_{21})^{(3)}(a_{22})^{(3)}(q_{20})^{(3)}G_{20}^{*}\right)$$

$$\left(\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(22)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(22)}T_{20}^{*}\right)\}=0$$

$$+$$

$$\left((\lambda)^{(4)}+(b_{26}')^{(4)}-(r_{26})^{(4)}\right)\{\left((\lambda)^{(4)}+(a_{26}')^{(4)}+(p_{26})^{(4)}\right)$$

$$\left[\left(\left((\lambda)^{(4)}+(a_{24}')^{(4)}+(p_{24})^{(4)}\right)(q_{25})^{(4)}G_{25}^{*}+(a_{25})^{(4)}(q_{24})^{(4)}G_{24}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(4)}+(b_{24}')^{(4)}-(r_{24})^{(4)}\right)s_{(25),(25)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(25)}T_{25}^{*}\right)$$

$$+\left(\left((\lambda)^{(4)}+(a_{25}')^{(4)}+(p_{25})^{(4)}\right)(q_{24})^{(4)}G_{24}^{*}+(a_{24})^{(4)}(q_{25})^{(4)}G_{25}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}+(b_{24}')^{(4)}-(r_{24})^{(4)}\right)s_{(25),(24)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(24)}T_{24}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}\right)^2+\left((a_{24}')^{(4)}+(a_{25}')^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)$$

$$\left(\left((\lambda)^{(4)}\right)^2+\left((b_{24}')^{(4)}+(b_{25}')^{(4)}-(r_{24})^{(4)}+(r_{25})^{(4)}\right)(\lambda)^{(4)}\right)$$

$$+\left(\left((\lambda)^{(4)}\right)^2+\left((a_{24}^{'})^{(4)}+(a_{25}^{'})^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)(q_{26})^{(4)}G_{26}$$

$$+\left((\lambda)^{(4)}+(a_{24}^{'})^{(4)}+(p_{24})^{(4)}\right)\left((a_{26})^{(4)}(q_{25})^{(4)}G_{25}^{*}+(a_{25})^{(4)}(a_{26})^{(4)}(q_{24})^{(4)}G_{24}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}+(b_{24}^{'})^{(4)}-(r_{24})^{(4)}\right)s_{(25),(26)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(26)}T_{24}^{*}\right)\}=0$$

$$+$$

$$(\lambda)^{(5)}+(b_{30}^{'})^{(5)}-(r_{30})^{(5)}\}\{\left((\lambda)^{(5)}+(a_{30}^{'})^{(5)}+(p_{30})^{(5)}\right)$$

$$\left[\left(\left((\lambda)^{(5)}+(a_{28}^{'})^{(5)}+(p_{28})^{(5)}\right)(q_{29})^{(5)}G_{29}^{*}+(a_{29})^{(5)}(q_{28})^{(5)}G_{28}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(5)}+(b_{28}^{'})^{(5)}-(r_{28})^{(5)}\right)s_{(29),(29)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(29)}T_{29}^{*}\right)$$

$$+\left(\left((\lambda)^{(5)}+(a_{29}^{'})^{(5)}+(p_{29})^{(5)}\right)(q_{28})^{(5)}G_{28}^{*}+(a_{28})^{(5)}(q_{29})^{(5)}G_{29}^{*}\right)$$

$$\left(\left((\lambda)^{(5)}+(b_{28}^{'})^{(5)}-(r_{28})^{(5)}\right)s_{(29),(28)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(28)}T_{28}^{*}\right)$$

$$\left(\left((\lambda)^{(5)}\right)^2+\left((a_{28}^{'})^{(5)}+(a_{29}^{'})^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)$$

$$\left(\left((\lambda)^{(5)}\right)^2+\left((b_{28}^{'})^{(5)}+(b_{29}^{'})^{(5)}-(r_{28})^{(5)}+(r_{29})^{(5)}\right)(\lambda)^{(5)}\right)$$

$$+\left(\left((\lambda)^{(5)}\right)^2+\left((a_{28}^{'})^{(5)}+(a_{29}^{'})^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)(q_{30})^{(5)}G_{30}$$

$$+\left((\lambda)^{(5)}+(a_{28}^{'})^{(5)}+(p_{28})^{(5)}\right)\left((a_{30})^{(5)}(q_{29})^{(5)}G_{29}^{*}+(a_{29})^{(5)}(a_{30})^{(5)}(q_{28})^{(5)}G_{28}^{*}\right)$$

$$\left(\left((\lambda)^{(5)}+(b_{28}^{'})^{(5)}-(r_{28})^{(5)}\right)s_{(29),(30)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(30)}T_{28}^{*}\right)\}=0$$

$$+$$

$$(\lambda)^{(6)}+(b_{34}^{'})^{(6)}-(r_{34})^{(6)}\}\{\left((\lambda)^{(6)}+(a_{34}^{'})^{(6)}+(p_{34})^{(6)}\right)$$

$$\left[\left(\left((\lambda)^{(6)}+(a_{32}^{'})^{(6)}+(p_{32})^{(6)}\right)(q_{33})^{(6)}G_{33}^{*}+(a_{33})^{(6)}(q_{32})^{(6)}G_{32}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(6)}+(b_{32}^{'})^{(6)}-(r_{32})^{(6)}\right)s_{(33),(33)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(33)}T_{33}^{*}\right)$$

$$+\left(\left((\lambda)^{(6)}+(a_{33}^{'})^{(6)}+(p_{33})^{(6)}\right)(q_{32})^{(6)}G_{32}^{*}+(a_{32})^{(6)}(q_{33})^{(6)}G_{33}^{*}\right)$$

$$\left(\left((\lambda)^{(6)}+(b_{32}^{'})^{(6)}-(r_{32})^{(6)}\right)s_{(33),(32)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(32)}T_{32}^{*}\right)$$

$$\left(\left((\lambda)^{(6)}\right)^2+\left((a_{32}^{'})^{(6)}+(a_{33}^{'})^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\right)$$

$$\left(\left((\lambda)^{(6)}\right)^2+\left((b_{32}^{'})^{(6)}+(b_{33}^{'})^{(6)}-(r_{32})^{(6)}+(r_{33})^{(6)}\right)(\lambda)^{(6)}\right)$$

$$+\left(\left((\lambda)^{(6)}\right)^2+\left((a'_{32})^{(6)}+(a'_{33})^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\right)(q_{34})^{(6)}G_{34}$$

$$+\left((\lambda)^{(6)}+(a'_{32})^{(6)}+(p_{32})^{(6)}\right)\left((a_{34})^{(6)}(q_{33})^{(6)}G^*_{33}+(a_{33})^{(6)}(a_{34})^{(6)}(q_{32})^{(6)}G^*_{32}\right)$$

$$\left((\lambda)^{(6)}+(b'_{32})^{(6)}-(r_{32})^{(6)}\right)s_{(33),(34)}T^*_{33}+(b_{33})^{(6)}s_{(32),(34)}T^*_{32}\Big)\}=0$$

$$+$$

$$(\lambda)^{(7)}+(b'_{38})^{(7)}-(r_{38})^{(7)}\}\{\left((\lambda)^{(7)}+(a'_{38})^{(7)}+(p_{38})^{(7)}\right)$$

$$\left[\left(\left((\lambda)^{(7)}+(a'_{36})^{(7)}+(p_{36})^{(7)}\right)(q_{37})^{(7)}G^*_{37}+(a_{37})^{(7)}(q_{36})^{(7)}G^*_{36}\right)\right]$$

$$\left(\left((\lambda)^{(7)}+(b'_{36})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(37)}T^*_{37}+(b_{37})^{(7)}s_{(36),(37)}T^*_{37}\right)$$

$$+\left(\left((\lambda)^{(7)}+(a'_{37})^{(7)}+(p_{37})^{(7)}\right)(q_{36})^{(7)}G^*_{36}+(a_{36})^{(7)}(q_{37})^{(7)}G^*_{37}\right)$$

$$\left(\left((\lambda)^{(7)}+(b'_{36})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(36)}T^*_{37}+(b_{37})^{(7)}s_{(36),(36)}T^*_{36}\right)$$

$$\left(\left((\lambda)^{(7)}\right)^2+\left((a'_{36})^{(7)}+(a'_{37})^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\right)$$

$$\left(\left((\lambda)^{(7)}\right)^2+\left((b'_{36})^{(7)}+(b'_{37})^{(7)}-(r_{36})^{(7)}+(r_{37})^{(7)}\right)(\lambda)^{(7)}\right)$$

$$+\left(\left((\lambda)^{(7)}\right)^2+\left((a'_{36})^{(7)}+(a'_{37})^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\right)(q_{38})^{(7)}G_{38}$$

$$+\left((\lambda)^{(7)}+(a'_{36})^{(7)}+(p_{36})^{(7)}\right)\left((a_{38})^{(7)}(q_{37})^{(7)}G^*_{37}+(a_{37})^{(7)}(a_{38})^{(7)}(q_{36})^{(7)}G^*_{36}\right)$$

$$\left(\left((\lambda)^{(7)}+(b'_{36})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(38)}T^*_{37}+(b_{37})^{(7)}s_{(36),(38)}T^*_{36}\right)\}=0$$

$$+$$

$$(\lambda)^{(8)}+(b'_{42})^{(8)}-(r_{42})^{(8)}\}\{\left((\lambda)^{(8)}+(a'_{42})^{(8)}+(p_{42})^{(8)}\right)$$

$$\left[\left(\left((\lambda)^{(8)}+(a'_{40})^{(8)}+(p_{40})^{(8)}\right)(q_{41})^{(8)}G^*_{41}+(a_{41})^{(8)}(q_{40})^{(8)}G^*_{40}\right)\right]$$

$$\left(\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(41)}T^*_{41}+(b_{41})^{(8)}s_{(40),(41)}T^*_{41}\right)$$

$$+\left(\left((\lambda)^{(8)}+(a'_{41})^{(8)}+(p_{41})^{(8)}\right)(q_{40})^{(8)}G^*_{40}+(a_{40})^{(8)}(q_{41})^{(8)}G^*_{41}\right)$$

$$\left(\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(40)}T^*_{41}+(b_{41})^{(8)}s_{(40),(40)}T^*_{40}\right)$$

$$\left(\left((\lambda)^{(8)}\right)^2+\left((a'_{40})^{(8)}+(a'_{41})^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\right)$$

$$\left((\lambda)^{(8)}\right)^2+\left((b'_{40})^{(8)}+(b'_{41})^{(8)}-(r_{40})^{(8)}+(r_{41})^{(8)}\right)(\lambda)^{(8)}\Big)$$

$$+\left(\left((\lambda)^{(8)}\right)^2+\left((a'_{40})^{(8)}+(a'_{41})^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\right)(q_{42})^{(8)}G_{42}$$

$$+\left((\lambda)^{(8)}+(a'_{40})^{(8)}+(p_{40})^{(8)}\right)\left((a_{42})^{(8)}(q_{41})^{(8)}G^*_{41}+(a_{41})^{(8)}(a_{42})^{(8)}(q_{40})^{(8)}G^*_{40}\right)$$

$$\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(42)}T^*_{41}+(b_{41})^{(8)}s_{(40),(42)}T^*_{40}\Big)\}=0$$

$$+$$

$$(\lambda)^{(9)}+(b'_{46})^{(9)}-(r_{46})^{(9)}\}\{\left((\lambda)^{(9)}+(a'_{46})^{(9)}+(p_{46})^{(9)}\right)$$

$$\left[\left(\left((\lambda)^{(9)}+(a'_{44})^{(9)}+(p_{44})^{(9)}\right)(q_{45})^{(9)}G^*_{45}+(a_{45})^{(9)}(q_{44})^{(9)}G^*_{44}\right)\right]$$

$$\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(45)}T^*_{45}+(b_{45})^{(9)}s_{(44),(45)}T^*_{45}\Big)$$

$$+\left(\left((\lambda)^{(9)}+(a'_{45})^{(9)}+(p_{45})^{(9)}\right)(q_{44})^{(9)}G^*_{44}+(a_{44})^{(9)}(q_{45})^{(9)}G^*_{45}\right)$$

$$\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(44)}T^*_{45}+(b_{45})^{(9)}s_{(44),(44)}T^*_{44}\Big)$$

$$\left((\lambda)^{(9)}\right)^2+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\Big)$$

$$\left((\lambda)^{(9)}\right)^2+\left((b'_{44})^{(9)}+(b'_{45})^{(9)}-(r_{44})^{(9)}+(r_{45})^{(9)}\right)(\lambda)^{(9)}\Big)$$

$$+\left(\left((\lambda)^{(9)}\right)^2+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\right)(q_{46})^{(9)}G_{46}$$

$$+\left((\lambda)^{(9)}+(a'_{44})^{(9)}+(p_{44})^{(9)}\right)\left((a_{46})^{(9)}(q_{45})^{(9)}G^*_{45}+(a_{45})^{(9)}(a_{46})^{(9)}(q_{44})^{(9)}G^*_{44}\right)$$

$$\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(46)}T^*_{45}+(b_{45})^{(9)}s_{(44),(46)}T^*_{44}\Big)\}=0$$

**And as one sees, all the coefficients are positive. It follows that all the roots have negative real part, and this proves the theorem.**

Note: **Same dovetailing explanation holds good for d/dt (partial differential with respect to t), d/dt,d^2/dt^2 (acceralation: double dot). Such an exposition is helpful in optimisation problems, duality of motion, wave motion a homogeneous functions studies, and control theory which we intend to incorporate in future.** It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the

concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

# Persistence, Coexistence And Thrusty Permanence For Interacting Structured Particles, Transactions In Finance, Econometric Productivity Functions, Environment, Interfaces, Cosmological Interactions Et Al

## Josef Hofbauer And Sebastian J. Schreiber Model Thereof: A Functional Appraisal

ABSTRACT. **JOSEF HOFBAUER AND SEBASTIAN J. SCHREIBER** studied the dynamics of interacting structured systems (in their case population dynamics) can be modeled by $\frac{dxi}{dt} = A_i\ (x)\ xi$ where $x_i\ \epsilon\ R^n_{i,}\ x = (x_1, \ldots, x_k)$, and $A_i(x)$ are matrices with non-negative off- diagonal entries. These models are permanent if there exists a positive global attractor and are robustly permanent if they remain permanent following perturbations of $A_i\ (x)$. Necessary and sufficient conditions for robust permanence are derived using dominant Lyapunov exponents $X_i\ (\mu)$ of the $A_i\ (x)$ with respect to invariant measures $\mu$. The necessary conditions requires $\max_i\ X_i(\mu)>0$ for all ergodic measures with support in the boundary of the non-negative cone. The sufficient condition requires that the boundary admits a Morse decomposition such that $\max_i$ Xi$(\mu)>0$ for all invariant measures $\mu$ supported by a component of the Morse decomposition. When the Morse components are Axiom A, uniquely ergodic, or support all but one population, the necessary and sufficient conditions are equivalent. Applications to spatial ecology, epidemiology, and gene networks are given. Appeared in *Journal of Differential Equations*, 248, 1955-1971 (2010) It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## INTRODUCTION

A fundamental issue in population biology is what are the minimal conditions to ensure the long-term survivorship for interacting populations whether they be viral particles, bio-chemicals, plants, or animals. When these conditions are met the interacting populations are said to persist or coexist. Since the pioneering work of Lotka and Volterra on competitive and predator prey interactions, Thompson, Nicholson, and Bailey on host-parasite interactions, and

Kermack and McKendrick on disease outbreaks, nonlinear difference and differential equations have been used to understand conditions for population persistence [1-5]. One particularly important form of persistence for deterministic models is *permanence* or *uniform persistence* which corresponds to the existence of a global attractor bounded away from extinction of one or more species. When such an attractor exists, interacting populations are able to recover from "vigorous shake ups" of the of the population state. Permanence has been characterized from a topological perspective [6-8] and with average Lyapunov functions [9-12] for differential equation models of the form

$$\frac{dxi}{dt} = x_i f_i(x) \qquad i = 1, \ldots, k \tag{3}$$

Where $x_i$ is the abundance of population $i$, $x = (x_1, \ldots, x_k)$, and $f_i(x)$ is the per-capita growth of population $i$.

An important extension of the concept of the permanence is *robust permanence* [13-16] which requires that (1) remain permanent after sufficiently small perturbations of the per-capita growth rates $f_i$. The importance of robust permanence as one crude form of structural stability for population models are approximations to reality. Consequently, if nearby models (e.g. more realistic models) are not permanent despite the focal model being permanent, and then one can draw few conclusions about the persistence of the biological system being approximated by the model. One can view robust permanence as one crude form of structural stability for debits-credits particles, propositions, transactions in the celestial subterranean realm and ceratoid dualism, finance like debiting say a current account and crediting to another account, or recrediting to the account which had been earlier debited, further conduction of the cash credit account, black hole formation, guzzling of inventories for the output, creation of income from output, reinvestment of income to investment,gormadizing of matter by black holes, and emission of radiation therefore, gamma ray bursts, in fact most of the transactions and interactions in the universe are conservative and quantum mechanical universe, something like a account acting as predator and prey. More often than not as we have seen from Haimovici model such models could be applied to the problems, propositions, which are not subsumed by the concomitant problems in the bearing. the population models. For this perspective, it is not unexpected that there are permanent systems that cannot be approximated by robustly permanent systems [17]. In [14, 16], criteria for robust permanence were developed with respect to the average per-capita growth rates $\int f_i(x)d\mu(x)$ with respect to invariant probability measures $\mu$ supported on the boundary of the positive orthant (i.e.

where one or more Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates are extinct). Roughly, these criteria for robust permanence require that the average per-capita growth rate (i.e. invasion rate) is positive for some missing species.

While (1) can account for many types of Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates. Interactions, it assumes that all individual within populations are exactly the same. However, theoretical biologists have long recognized that different individuals within a population may be in different states (e.g. different sizes or ages, living in different spatial locations) and these differences can have important consequences for population dynamics [18-24]. To account for how these differences influence persistence of interacting populations, we develop criteria for robust permanence for the dynamics of $k$ interacting structured populations. Roughly, these dynamics correspond to replacing $x_i$ in (1) by vectors and $f_i(x)$ by matrices. In section 2, we describe these models in greater detail and review some basic concepts from dynamical systems theory. In section 3, we introduce the structured analog of average per-capita growth rates and define the concept of robustly unsaturated invariant sets. Necessary and sufficient conditions for an invariant set to be robustly unsaturated are proven. For Axiom A or uniquely ergodic invariant sets, the necessary and sufficient criteria are shown to be equivalent. In section 4, we use Morse decompositions and the criteria for robustly unsaturated invariant sets to develop necessary and sufficient conditions for robust permanence. In sections 5 and 6, we provide several applications of our results to spatially structured ecological models, structured epidemiological models, and models of gene networks. We apply the model to econometrics. To give an example investment needs inventory to be used towards the consummation of output, part of output is to be sold for consolidation of income, and some income might be replenished in the investment and so on. In finance a cash credit account could be debited or credited when there is say purchase of demand bills or discount of usance bills. Cash credit account is debited whenever cheques are issued towards the purchase of raw materials and so on. In fact some of the models have been applied to Organic psychiatry and we have shown how such models are useful in dealing with the propositions of Husserl, Derrida, Deleuze and others.

## 3. MODELS AND ASSUMPTIONS

Let $x_i$ denote the state of the $i$-th population that lies in the non-negative cone $C_i$ of $R^{ni.}$ Define $x=(x_1,\ldots x_k)\epsilon C$ to be the non-negative cone of $R^n$ where $n=\sum_{i=1}^{k} n_i$. Let $C_i^+$ denote positive cone { $x_i \epsilon C_i$: $\prod_i$i noitalupop rof$^j$

$i$ *and* $C^+=\prod_i C^+_i$ , the positive cone for the interacting populations. If $x \mapsto A_i(x)$ is a map into $n_i x n_i$ matrices that describes the growth of population I, then the dynamics of the interacting populations are given by

$$(2)\frac{dx_i}{dt} = A_i(x)x_i \quad i=1,\ldots,k$$

Let $x.t$ denote the solution of (2), with the initial condition x, and more generally the semi flow generated by (2). Before stating our assumption on (2), recall a few definitions from dynamical systems. Given sets $I \subseteq R$ and $K \subseteq C$, let $K.I = \{x.t : t \epsilon I, x \epsilon K\}$. A set $K \subseteq C$ *is invariant* if K.t = K for all t>0, and *forward invariant* if $K.t \subset K$ for all $t > 0$. The *omega limit* set of a set $K \subseteq C$ equals $\omega(K) = \bigcap_{t\geq 0,} \overline{K.[t,\infty)}$. The *alpha limit* set of $K \subseteq C$ is $a(K) = \bigcap_{t\leq 0} \overline{K.(-\infty,t]}$. Given a forward invariant set K, $B \subset K$ is an attractor for the semi flow $x.t$ *restricted* to K provided there exists an open neighborhood $U \subseteq K$ of B such that $\omega(U) = B$. The stable set of a *compact invariant* set $K$ is defined by

$$W^s(K) = \{x \epsilon C : \omega(x) \neq \emptyset \text{ and } \omega(x) \subseteq K\}.$$

The semi flow generated by (2) is *dissipative* if there exists an attractor B with $W^s(B) = C$.

Throughout this paper, we make the following assumptions:

**A1**: $x \mapsto A_i(x)$ is continuous

**A2:** $A_i(x)$ is irreducible and has non-negative off- diagonal entries,

**A3:** $x.t$ is defined for all $t \geq 0$

**A4:** $x.t$ is dissipative with global attractor $\Gamma(A)$.

A1 is a basic regularity assumption. The non-negative of off- diagonal entries in A2 implies that there are no negative feedbacks between individuals of different states in population $i$. This assumption is meet for many types of structured models as discussed in sections 5 and 6. The irreducibility assumption A2 is generally meet and implies that all individuals within a population can pass through all states. In Remark1, we discuss how this irreducibility assumption can be relaxed. A3 ensures that the population dynamics are defined for all future time. A4 requires that

population densities/ abundances eventually are uniformly bounded, a condition that should be meet for an biologically realistic model.

We say that the semi flow of (2) is *permanent* if this semi flow is dissipative and there is a positive attractor $B \subset C^+$ such that $W^s(B) \subset C^+$. We will denote this positive attractor by $\Lambda(A)$. Permanence ensures that populations can recover from rare large perturbations and allows for a diversity of dynamical behaviors.

For K ⊂ C and $\delta > 0$, define the $\delta$- *neighborhood* of K as

$$N\delta(K) = \{x \in C : \backslash x - y < \delta \text{ for some } y \in K\}.$$

We define a $\delta$- *perturbation of (2)* to be a system of the form

$$\frac{dx}{dt} = \tilde{A}(x)x$$

that satisfies assumptions A1-A4, $\Gamma(\tilde{A}) \subset N$ *( $\Gamma(A)$), and* $||A(x) - \tilde{A}(x)|| \leq \delta$ *for all* $x \in N$ *($\Gamma(A)$).* (2) is *robustly permanent* if there exist >0 and $\epsilon > 0$ such that all $\delta$- perturbations of (2) are permanent, and $d(\Lambda(\tilde{A}), C\backslash C^+) > \epsilon$ for all $\delta$- perturbations i.e., there is a common/ uniform region of repulsion around the boundary.

## INUNDATION OF COMPACT SETS

We begin by studying the linear skew product flows on $\Gamma(A) \times R^n_i$ defined by $(x,y).t = (x.t, B_i(t,x)y)$ where $Y(t) = B_i(t,x)$ is the solution to $\frac{dY}{dt}(t) = A_i(x.t)Y(t)$ with y(0) equal to the identify matrix. Our assumption that $A_i$ is irreducible with non- negative off diagonal entries implies that $B_i$ (t,x)$C_i \subset C_i^+$ for all $x$ and t > 0 (see, e.g., [25] ). A result of Ruelle [26, Prop, 3.2] implies that there exist continuous maps $u_i\, v_i : \Gamma(A) \rightarrow C_i^+$ with $u_i(x)| = |v_i(x)| = 1$ such that The line bundle $E_i(x)$ spanned by $u_i(x)$ is invariant i.e. $E_i(x.t) = B_i(x,t)E_i(x)$ for all t ≥ 0.The vector bundle $F_i(x)$ perpendicular to $v_i(x)$ is invariant i.e. $F_i(x.t) = B_i(x,t)F_i(x)$ for all $t \geq 0$.There exist constants $a > 0$ *and* $\beta > 0$ such that(3) $||B_i(t,x)F_i(x)|| \leq a\, exp(-\beta t)||B_i(t,x)E_i(x)||$ For all $x \in K$ *and* $t \geq 0$.**Interaction rates.** Given any x $\in \Gamma$(A), define the *interaction and inundation rate* of *species i at particle/proposition/problem/ state x as* $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right)$ *in* $||B_i(t,x)||$ when $x_i = 0$, $\lambda_i(x)$ provides an upper bound to the rate of growth of population *i* when introduced at infinitesimally small densities. Important properties of this invasion rate are summarized in the following two propositions. For instance, the second proposition implies if $(x.t)_i$ stays bounded away from zero (i.e. population *i* persist), then

$\lambda_i(x)=0$. Alternatively, if $\lambda_i(x)<0$,then population $i$ is doomed to extinction. **Proposition 1.** $\lambda_i(x)$ *satisfies the following properties:* $\lambda_i(x)= limsup_{t\to\infty}\left(\frac{1}{t}\right) In \ |B_i \ (t,x)v|$ *for any* $v>0$ *in* $R^{ni}$, *and* $\lambda_i(x)= limsup_{t\to\infty}\left(\frac{1}{t}\right) \int^{t}_{0}(A_i(x.s)u_i(x.s),u_i(x.s))ds.$

**Proof:** To prove the first property, we first show that

$$\lambda i(x) = \underset{t\to\infty}{limsup}\frac{1}{t} In\ |B_i(t,x)u_i(x)|.$$

To this end, let $v \epsilon R^{ni}$ be any non-zero vector. Since $R^{ni}= E_{i(x)} \ F_i(x)$, we can write $v=au_i(x)+w$ with $a \epsilon R$ and $w \epsilon F_i(x)$. Equation (3) implies

$|B_i(t,x)v| \ \leq \ a|\ B_i(t,x)u_i(x)|+ |B_i(t,x)w| \ \leq \ |\ B_i(t,x)u_i(x)|(a+ \alpha+ exp\ (-\beta t)|\ w|)$. Hence, $\underset{t\to\infty}{limsup}\frac{1}{t} In\ B_i(t,x)\ v| \leq$

$$\underset{t\to\infty}{limsup}\frac{1}{t} In\ B_i\ (t,x)\ u_i\ (x) \leq \lambda_i(x)$$

for all non zero vectors $v \epsilon R^{ni}$. Since $||B_i(t,x)|| = sup|v|=1B_i(t,x)v|$, this inequality implies that $\lambda i(x) = \underset{t\to\infty}{\text{limsup}}\frac{1}{t}$ In $|B_i(t,x)v|$, this inequality implies that $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right)$ In $|t,x)u_i(x)|$. Now let v $\epsilon$ $C_i^{+}$. Then, we can write $v = au_i(x)+ w$ with $a >0$ and $w \epsilon F_i$.

Equation (3) implies

$$|B_i(t,x)v| \ \geq \ a|B_i(t,x)u_i(x)|-|B_i(t,x)w|$$

$$\geq |\ B_i(t,x)u_i(x)|(a- \alpha+ exp\ (-\beta t)|\ w|).$$

Since a>0, this inequality implies that

$$\lambda i(x) \geq \underset{t\to\infty}{\text{limsup}}\frac{1}{t} \text{In}\ |B_i(t,x)v| \ \geq \ \lambda i(x) \geq \underset{t\to\infty}{\text{limsup}}\frac{1}{t} \text{In}\ |B_i(t,x)u_i(x)| = \lambda i(x)$$

To prove the second assertion, let $b_i(t,x) = In\ |B_i(t,x)u_i(x)|$. Invariance of $u_i(x)$ implies that $b_i(t,x)$ is additive:

$$B_i(t+s,x) \ = \ In|B_i(t+s,x)\ u_i(x)|= In\ |B_i(s,x.t)B_i(t,x)u_i(x)$$

$$= \text{In}\ |B_i(s,x.t)\ u_i(x.t)||\ B_i(t,x)\ u_i(x)|$$

$$= \text{In}\ b_i(s,t.x)\ b_i(t,x)$$

Additivity of $b_i(t,x)$ and the fact that $b_i(0,x)$ =0 implies

$$\frac{d}{dt}b_i(t,x) = \lim_{s\to 0}\frac{bi(t+s,x-bi(t,x)}{S-}$$

$$= \lim_{s\to 0}\frac{bi(s,x.t)}{S} = \frac{d}{ds}|_{s=0}\, b_i(s,x.t)$$

=

$$=(A_i(x.t)u_i(x.s),u_i(x.s),$$

The Fundamental Theorem of Calculus implies

$$B_i(t,x)= \int_0^t (A_i(x.s)\, u_i(x.s),\, u_i(x.s))ds\backslash$$

and the second assertion follows.

**Proposition 2:** *For the solutions of (2) we have:* if $x_i>0$, then $\lambda i(x) \le 0$. If $\lambda i(x)<0$, then $\lim_{t\to\infty} (x.t)_i=0$ If $limsup_{t\to\infty}$ $(x.t)_i|>0$, then $\lambda i(x)=0$.

**Proof:** First, assume that x satisfies $x_i>0$. Since the semi flow is dissipative, there exists $\gamma >0$ such that $|x.t| \le \gamma$ for all $t \ge 0$. Proposition 1 and the definition of the skew product flow imply

$$\lambda i(x) = \operatorname{limsup}_{t\to\infty}\frac{1}{t}\,\mathrm{In}\,|B_i(t,x)x_i| = \operatorname{limsup}_{t\to\infty}\frac{1}{t}\,\mathrm{In}\,|\,(x.t)_i\,|$$

$$\le \operatorname{limsup}_{t\to\infty}\frac{In\,\gamma}{t} = 0$$

To prove the second assertion, assume that $\lambda i(x) < 0. If\ x_i=0$, then the invariance of the faces of C imply that $(x.t)_i=0$ for all $t\ge 0$. Alternatively, if $x_i> 0$, then Proposition I and the skew product flow imply

$$\operatorname{limsup}_{t\to\infty}\frac{1}{t}In|(x.t)_i| = \lambda i(x) < 0.$$

The final assertion follows from the first two assertions. 5

**3.2. Invariant measures.** We review some definitions from ergodic theory. Given a Borel probability measure $\mu$ on ***C***, the *support* of $\mu$, denoted supp( $\mu$),is the smallest closed set whose complement has measure zero. A Borel probability measure $\mu$ is called *invariant* for (2) provided that Integral $h(x)d\mu(x) =$ Integral $h(x.t)d\,\mu(x)$for all

$t \geq 0$ and for all bounded continuous functions $h$:C→R. An invariant measure $\mu$ is called *ergodic* provided that $\mu(B)$=0 or 1 for any invariant Borel set *B.*

For an invariant measure $\mu$ define the invasion rate species $i$ with respect to $\mu$ as

$$\lambda_i(\mu) = \int_c \langle A_i(x)u_i(x), u_i(x)\rangle d\mu(x)$$

**Proposition 3:** Let $\mu$ be an *invariant measure.* $\lambda_i(\mu)$ *Satisfies the following properties:* $\lambda_i(x) = \lim_{t\to\infty} \frac{1}{t}\int_0^t \langle A_i(x.s)u_i(x.s), u_i(x.s)\rangle, ds$ exists for $\mu$ - *almost every* $x$. *Moreover, if* $\mu$ *is ergodic,* $\lambda_i(x) = \lambda_i(\mu)$ for $\mu$-*almost every* $x$.*If* $\mu$ *is ergodic, then there exists* $I \subset \{1, \dots, k\}$such that $\mu(\Pi_{i\in I} C_i^+)$=1 and $\lambda_i(\mu) = 0$ *for all* $i \in I$.

**Proof:** Let $\mu$ be an invariant measure. Define $h_i(x) = \langle A_i(x)u_i(x), u_i(x)\rangle$. The Birkhoff's' ergodic theorem and Proposition 1 imply

$$\lambda_i(x) = \lim_{t\to\infty} \frac{1}{t}\int_0^t h_i(x.s)ds$$

Exists for $\mu$- almost every $x$. Moreover, $\lambda_i(x) = \lambda_i(\mu)$ $\mu$- almost surely if $\mu$ is ergodic.

Assume $\mu$ is ergodic. By ergodicity and invariance of the faces of **C**, there exists $I \subset \{1, \dots, k\}$ such that $\mu(\Pi_{i\in I} C_i^+)$=1. Let K $\subset \Pi_{i\in I} C_i^+$ be a compact set such that $\mu(K) > 0$. The Poincare recurrence theorem and the Birkhoff ergodic theorem imply that there is $x \in \Pi_{i\in I} C_i^+$such that $x.t_n \in K$ for some $t_n \uparrow \infty$ and $\lambda_i(x) = \lambda_i(\mu)$ for all $i$. The third assertion of Proposition 2 implies $\lambda_i(\mu) = 0$ for all $i \in I$.

**3.3. Robustly unsaturated sets**. Let K $\subset \partial C$ be a compact isolated invariant set for the flow $x.t$ restricted to $\partial C$.*K* is unsaturated if Ws (K) $\subset \partial C$ and K is isolated for $x.t$. If *K* is not unsaturated, then *K* is saturated. *K* is robustly unsaturated for (2) if there exists $\delta > 0$ such that the continuation of K for any $\delta -$ perturbation of (2) is unsaturated.**Theorem 1:** Let *K* be a compact isolated invariant set for $x.t$ restricted to$\partial C$. If one of the following equivalent conditions hold For all invariant measures $\mu$ supported by *K* (4) $\max_{1\leq i\leq k} \lambda_i(\mu) = 0$

There exist $p_i > 0$ such that (5) $\sum_{1\leq i\leq k} p_i \lambda_i(\mu) = 0$

For all ergodic probability measures supported by *K*.

then $K$ is robustly unsaturated for (2). Alternatively, if $x \to A(x)$ is twice continuously differentiable and $K$ is robustly unsaturated, then (4) holds for all ergodic measures $\mu$ supported by *K*.

**Remark 1**: For some applications (e.g., the disease model considered in section 6.1) it useful to relax the irreducibility assumption A2. For instance, if there exists an open neighborhood U of K such that for each $i$, $A_i(x)$ has a fixed off-diagonal sign pattern for all $x \in U$,then $A_i(x)$ can be decomposed into a finite number, say $m_i$, of irreducible components. For each of these irreducible components, one can define $\lambda_i^j(\mu) = \int_C \langle A_i^j(x) u_i^j(x), u_i^j(x) \rangle \, d\mu(x)$where $A_i^j(x)$ is the submatrix of $A_i(x)$ corresponding to the *j*-th irreducible component of $A_i(x)$ *and* $u_i^j(x)$ is the continuous invariant sub bundle for the irreducible linear cocycle determined by $A_i^j(x)$. If we define $\lambda_i(\mu) = \max_{1 \le j \le m_i} \lambda_i^j(\mu)$, then all of the assertions of Theorem 1 still hold.

**Proof:** To prove the first assertion of the theorem, we need the following lemma. We call an invariant measure $\mu$ for (2) saturated if $\lambda_i(\mu) \le 0$ for all $i$

**Lemma 1**: Let $\delta_n$ be a non –negative sequence that converges to zero as n$\to \infty$. If $\mu_n$ are saturated invariant measures for $\delta_n -$ perturbations of (2), then the weak$^*$ limit points of $\{\mu_n\}_{n=1}^{\infty}$ is a non-empty set consisting of saturated invariant measures for (2).

**Proof:** Let $\delta_n$ be a non-negative sequence that coverages to zero. Let $\mu_n$ be saturated invariant measures for $\delta_n$ perturbations, $\dot{x} = A^n(x)x$, of $\dot{x} = A(x)x$. By weak$^*$ compactness of Borel probability measures supported on $N_1(\Gamma(A))$, there exist weak$^*$ limit points of $\mu_n$ . Let $\mu$ be such a weak$^*$ limit point. Since $\mu_n$ are supported by $N_1(\Gamma(A))$, for all $n$ sufficiently large, $\mu$ is supported by $N_1(\Gamma(A))$.To verify that $\mu$ is an invariant measure for (2),let $h: C \to R$ be a bounded continuous function. Let $t > 0$ and $\in > 0$ be given. Choose $n$ sufficiently large so that $|\int_c h(x)\, d\mu(x) - \int_c h(x) d\mu_n(x)| \le \in$, $|\int_c h(x.t)\, d\mu(x) - \int_c h(x.t) d\mu_n(x)| \le \in$, and $|h(x.t) - h(x^n.t)| \le \in$ for all $x \in N_1(\Gamma(A))$.Then $|\int_c h(x.t) - h(x) d\mu(x)|$ is

$\leq \quad |\int_c h(x.t)\,d\mu(x) - \int_c h(x.t)\,d\mu_n(x)| + |\int_c h(x.t)\,d\mu_n(x) - \int_c h(x)\,d\mu(x)|$

$\leq \quad \in + \int_c |h(x.t) - h(x^n.t)|\,d\mu_n(x)| + |\int_c h(x^n.t)|\,d\mu_n(x) - \int_c h(x)d\mu(x)|$

$\leq \quad 2 \in + |\int_c h(x^n.t) - h(x)\,d\mu_n(x)| + |\int_c h(x)d\mu_n(x) - \int_c h(x)d\mu(x)|$

$\leq \quad 3 \in$

Where the last line follows from the invariance of $\mu_n$ for $x^n.t$. Since $\in > 0$ and t > 0 are arbitrary, we have $\int_c h(x.t)d\mu(x) = \int_c h(x)d\mu(x)$ for all t >0 and all bounded continuous functions $h: C \to R$. It follows that $\mu$ is an invariant measure for (2). To see that $\mu$ is saturated, define $h_i^n(x) = \langle A_i^n(x)u_i^n(x), u_i^n(x)\rangle$ where $u_i^n(x)$ spans the invariant one-dimensional bundle given by $[26, Prop.3.2]$ for $\frac{dx}{dt} = A^n(x)x$. Since $A_i^n(x) \to A_i(x)$ uniformly for $x \in N_1((A))$ as $n \to \infty$, $[26, Prop.3.2]$ implies that $u_i^n(x) \to u_i(x)$ converges uniformly for $x \in N_1(\Lambda(A))$. Given $\in > 0$, choose $n$ sufficiently large so that $|h_i^n(x) - h_i(x)| \leq \in$ for all $x \in N_1(\Lambda(A))$ and $|\int_c h_i^n(x)\,d\mu(x) - \int_c h_i^n(x)\,d\mu_n(x)| \leq \in$.

Then

$\lambda_i(\mu) \leq \int_c |h_i(x) - h_i^n(x)|d\mu(x) + |\int_c h_i^n(x)\,d\mu(x) - \int_c h_i^n(x)\,d\mu_n(x)|$

$+ \int_c h_i^n(x)\,d\mu_n(x)$

$\leq 2\in$

Since $\in > 0$ is arbitrary, $\lambda_i(\mu) \leq 0$ for all $i$ and $\mu$ is saturated.

Using this lemma, we prove that if K is saturated, then there exists a saturated invariant measure $\mu$ supported by K. Assume K is saturated. Work of Hofbaue and So $[7, Thm.2.1]$ implies that either there exists $y \in W^s(K) \cap C^+$ or K is not isolated for the unrestricted flow $x.t$ If there exists $y \in W^s(K) \cap C^+$, then for all $t \leq 0$ define $v_t = \frac{1}{t}\int_0^t \delta_{y.s}\,ds$ where $\delta_{y.s}$ denotes a Dirac measure based at the point $y.s$. Dissipativeness of (2) and weak* compactness of the Borel probability measures supported by (A) imply there exist $t_k \to \infty$ such that the sequence $v_{tk}$ converges in the weak* topology to a Borel probability measure $\mu$ with support in K.A standard argument

implies that $\mu$ is $x.t$ invariant. Define $h_i$: C → R by $h_i(x) = (A_i(x)u_i(x), u_i(x))$.Proposition 2 and weak* convergence imply that

$$0 \geq (\lambda i(y) = \limsup_{t\to\infty} \frac{I}{t}\int_0^t h_i(y.s)ds$$

$$= \limsup_{t\to\infty} \int_c h_i \, dv_t$$

$$\geq \limsup_{t\to\infty} \int_p h_i \, d\, v_{tk} = \lambda i(\mu)$$

For all $i$. Alternatively suppose that K is not isolated for the semi flow Then there exists a sequence of positive $\omega - limit$ sets that accumulate on K. Let $\mu_n$ be a sequence of ergodic probability measures supported by these $\omega -$ limit sets. Propositions 1 and 3 imply that $\mu_n$ are saturated for all $n$ Applying Lemma I with $\delta_n$=0 for all $n$ implies that there exists a saturated invariant measure $\mu$ suppoorted by *K*. Hence we have shown that if *K* is saturated then there exists a saturated invariant probability measure $\mu$ supported by *K* Equivalently, (4) holding for all invariant probability measures with support in *K* implies that *K* is unsaturated.

Next we show that if *K* not robustly saturated, then there exists a saturated invariant measure supported by *K* Indeed suppose *K* is not robustly saturated, Then there exists a non-negative sequence $\delta_n$ converging to zero and a sequence of saturated measures $\mu_n$ for $\delta_n$-perturbations of (2)with support in the continuation of *K*. Let $\mu$ be a weak* limit point of $\{\mu_n\}\infty_{n=1}$ Lemma 1 implies that $\mu$ is saturated. Moreover, since the continuation of *K* converges to *K* as $\delta_n \to 0, \mu$ is a saturated invariant measure for (2) supported by *K* Hence we have shown that if *K* is not robustly saturated then there exists a saturated invariant measure supported by *K* Equivalently if (4) holds for all invariant measures with support in *K* then *K* is robustly unsaturated.

To see the equivalence of the conditions given by (4) and (5) let $\Delta = \{p \in R_+^k : \sum_i P_i = 1\}$ and notice that

$$\text{Min max } \lambda i\ (\mu = \min_{\mu} \max_{p\epsilon\Delta} \sum_i Pi\lambda i(\mu)$$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K* The Minimax theorem (see, e. g.,[27] implies that

(6) $\min \max \lambda i(\mu) = \min_{\mu} \max_{p \epsilon \Delta} \sum_{i} Pi\lambda i(\mu)$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K*. Since min$\mu$ $\sum_{i} P_i \lambda_i(\mu)$ is attained at an ergodic probability measure with support in *K*, the equivalence of the conditions given by (4) and (5) is established.

To prove the final assertion of the Theorem assume $x \to A(x)$ is twice continuously differentiable and there exists a saturated ergodic measure $\mu$ supported by *K* Proposition 3 implies that there exists *I*C {1,…,k} such that (□$i \in iC_i^+$)=1and $\lambda i(\mu) = 0$ For all $i \in I. SinceKc\,.\partial$□,$\{1, \ldots . k\}\backslash I$ is non-empty. We will show that (2) is not robustly permanent by proving that for all $\delta > 0$ there exists a $\delta$ −perturbation of (2) for which the continuation of *K* is saturated Let $\delta > 0$ be given Choose V $c$ N$\delta$ (*Γ*(A)) to be a compact neighborhood of Γ(A) such that V.t $c$ int *V* for all t > 0.Let $\eta > 0$ be such that *W*=:*N* $\eta$(Γ(A))$c$ *V* Let p:c→ $[0,1] be$ a$C^{\infty}$function such that $p(x) = 1 \quad for\ all\ x \in \Gamma(A) and\ p(x) = 0 for\ all\ x \in C\backslash W\ Define\ \tilde{A} = (\tilde{A}_{1,}\ldots.,\tilde{A}_{k})$ by

$$\tilde{A}_i(x) = \begin{cases} \frac{A(x)}{A(x) - \frac{\delta}{2Idp\,(x)}} & if\ i \epsilon I \\ if\ i \notin I \end{cases}$$

Where Id denotes the identity matrix of appropriate dimension Let x. t denote the semi flow of $x = \tilde{A}(x)x$.Since$x.t = x.t$ $whenever\ x.[0,t] \in C\backslash W, it\ follows\ that\ \ \tilde{V}.t \subset intV\ for\ all\ t > 0. Hence, \Gamma(\tilde{A})c\ V\ c\ N\delta\ (\ (A))). We\ \ also\ have\ ||A(x) - \tilde{A}(x)|| \leq \frac{\delta}{2}$Therefore $x = \tilde{A}(x) is\ a\ \delta -$ perturbation of (2) By construcation,$x.t = x.t\ for\ all\ x \in K and\ t \geq 0. Hence, \mu$ is ergodic for the semiflow of $\dot{x} = \tilde{A}(x)x\ and\ \lambda_i(\mu) \leq -\frac{\delta}{2}$ for this semiflow and $i \in \{1, \ldots, k\}\backslash I$.Let $\mathcal{L}$ and O be the Lyapunov exponents and Oseledec regular points supported by $\mu$ for $x = \tilde{A}(x)x(see, e.g.,$

[28] for definitions).At each point $x \in O$, the splitting of $R^n$ determines three subspaces: the stable subspace $E^s(x)$, the center subspace $E^c(x)$, and the unstable subspace $E^u(x)$ Proposition 3 and our choice of Ã imply that $E^s(x) \cap C^+ \neq \emptyset$. The Pesin stable manifold theorem [28, Corollaries 3.17 and 3.18] implies that tangent to $E^s(x)$ $E^c(x)$ and $E^u(x)$ are locally $x.t-$invariant families of $C^1$ discs $W^s_x, W^c_x$ and $W^u_x$ corresponding to the stable ,center and unstable manifolds The family of stable manifolds $W^s_x$ is contained in $W^s$ (Ã,supp $(\mu)$). $Since\ E^s(x) \cap C^+ \neq \emptyset$, $W^s_x \cap C^+ \neq \emptyset$ for some $x \in K$. Consequently $K$ is saturated for $x = \tilde{A}(x)x$.

The difference between the sufficient and necessary condition is that for the sufficient condition (4) has to hold for all invariant measures supported by *K*, while for the necessary condition (4) has to hold only for ergodic measures supported by *K*. since the invariant measures lie in the convex hull of the ergodic measures the sufficient condition can be more restrictive then the necessary condition.

For three Classes of invariant sets the necessary and sufficient conditions coincide Recall a compact invariant set *K* for (2) is uniquely ergodic if K only supports one invariant measure Recall a compact invariant set *K* is Axiom A if the flow of (2) restricted to *K* is transitive and hyperbolic (see, e.g.,[29] for definitions.)

**Corollary 1:** Assume $x \to A(x)$ is twice continuously differentiable It K$\subset \partial$□ is a compact invariant set and either is uniquely ergodic or Axiom A, then *K* is robustly unsaturated if and only if (4) holds for all ergodic, measures $\mu$ with support in *K*

Proof. If K is uniquely ergodic then the assertion follows immediately For an Axiom A invariant set Sigmund [30, thm.1] has proven that invariant measures supported by periodic orbits of *K* are dense in the set of invariant probability measures supported by *K* In particular any invariant measure can be approximated by an ergodic measure and consequently the result follows.

Another special case where the necessary and sufficient conditions coincide is when *K* supports all populations except one for discrete time models this case was considered by [22].

**Corollary2:** Assume $x \to A(x)$is twice continuously differentiable Let K$c$ $II_i \geq 2C_i^+$ be a compact invariant set. Then K is robustly unsaturated if and only if $\lambda_1(\mu) > 0$ for all invariant measures supported by *K*

## 4. MORES DECOMPOSITIONS AND ROBUST PERMANENCE

To state the sufficient condition for robust permanence, we use a characterization of permanence due to Garay [8] and Hofbauer and So [7] that involves Morse decompositions of the boundary flow Conley[31] defined a collection of sets$\{M_{1,\ldots,}M_m\}$ to be a Morse decomposition for a compact invariant set $K$ if $M_1\ldots,M_m$ are pairwise disjoint, compact isolated invariant sent for the flow of (2) restricted to K. For each $x \in K$ there are integers r=r$(x) \leq s = s(x)$ such that $a(x) \subseteq M_r$and $\omega(x) \subseteq M_{s.}$ If $r(x) = s(x), then\ x \in M_{r.}$ Garay, Hofbauer and so [7,8] proved the following characterization of permamnce

**Theorem 2** (Garay, Hofbauer-so) if $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$then (2) is permanent if and only if each of the components $M_1$ are unsaturated.

Theorem 1 and 2 imply the following result: As stated in the above all these results are true for most of problems, interactions, transactions in cosmology, econometrics, finance and other subjects, which we shall illustrate as we go on constructing models for the same.

**Theorem 3:** If $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$and (4) holds for each of the components of the More decomposition, then (2) is robustly permanent Conversely if $x \to A(x) is\ twice\ continuously\ differentiable$

*and (4) is violated by an ergodic measure supported by one of the components of the Morse decomposition, then (2) is not robustly permanent.* Theorem 3 in conjunction with Corollaries 1 and 2 yield a characterization of robust permanence for a class of structured models.

**Corollary 3:** *Let $\{M_1,\ldots..,M_m\}$ os a Morse decomposition for* $\partial$C. *Assume x* $\mapsto$ *A(x) is twise continuously differentiable and for each Morse component $M_i$ one of the following assertions hold $M_i$ is Axiom A, $M_i$ is uniquely ergodic, or There exists j* $\in$ { 1,….,k} *such that $M_i$ ----- { $x \in C : x_l \gg 0$ for $l \neq j$} Then (2) is robustly permanent if and only if (4) holds for all ergodic measures supported by $U_iM_i$.*

Our result also provided a structured analogue for characterizing totally permanent systems[32]. For an ergodic probability measure$\mu$, define sup ($\mu$) to be the subset $I$ -----{1,…,k} such that $\mu(---_i \in_I C_+) = 1$

**Corollary 4:** *The following statements are equivalent: For all ergodic probability measures μ with support in ∂C, $\lambda_i(\mu)>0$ for all $I \in \{1,\ldots,k\} \setminus supp(\mu)$ (2) and all of its subsystems are robustly permanent.*

*Proof. Suppose the first statement holds. Then unsaturated condition given by (5) with p= (1,…,1) and Theorem 3 implies that (2) and all its subsystems are robustly permanent.. The other direction follows form the necessary condition for robust permanence in Theorem 3 and the second assertion of Proposition 3.*

*To illustrate the broad applicability of theorem 3, we develop applications to spatially structured models in section 5 and to a disease model and a gene network model in section 6.*

## 5. APPLICATIONS: PATCH MODELS

A fundamental application of our result are to spatially structured models, with *k* species dispersing between *m* patches {33-35}:

(7) $$\frac{dx_i^j}{dt} = x_i^j f_i^j(x^j) + \sum_l d_i^{jl} x_i^l - e_i^j x_i^j \qquad i = 1,..,k \qquad j=1,\ldots,m$$

Where $x_i^j$ denotes the density of species I in patch *j*, $x^j = (x^j_1,.....,x^j_k)$ is the vector of species densities in patch *j*, $f^j_i$ is the per-capita growth rate of species *i* in patch *j*, $d_i^{jl} \geq 0$ is the dispersal rate for species *i* from patch *l* into patch *j*, and $e_i^j$ is the emigration rate of species *i* out of patch *j*. Hence $e_i^j \geq \sum_l d_i^{lj}$. We assume that the matrices $(d_i^{lj})_{j,l}$ are irreducible for each species *i* . it is then easy to write (7) in the form (2), with $x_i = (x^l_i, \ldots, x_i^m)$ 'where' denotes transpose.

For a single species (7) generates a monotone flow. Under mild assumption, e.g., each $f^j_i$ *is* decreasing and negative for large densities, there is a globally stable equilibrium, see, e.g.,[36] and [37, sect. 5.4]. This equilibrium will be positive if the invasion rate at the origin λ (0), which is given by the leading eigenvalue (stability modulus) of the matrix $(d^{jl} + f^j(0) - e^j)$, is positive.

5.1 Two types of debitable and creditable systems:. For *two competing species*, e.g., $f^j_i$ is decreasing with respect to both species and is negative at large densities, the dynamics are still monotone, and hence almost all orbits converge

to an equilibrium. Robust permanence requires two conditions. First, $\lambda_i$ (0) >0 for both species in which case there are two single species equilibria $E_1$ and $E_2$. Second, the invasion rates $\lambda_2$ ($E_1$ ) and $\lambda_2$ ($E_2$) are positive. Whether there is a unique positive globally stable equilibrium for the robustly permanent system depends in a delicate way on the system parameters, see [38]. For slowly dispersing populations (i.e. $0< e^j_i \ll 1$ for all *i, j* ), the robust permanence condition is particularly

Straight forward to verify. Let $E^j_i$ be the largest solution to $x^j_i \quad f_i{}^j$ ($x^j_i$) = 0 i.e. the equilibrium attained by species *I* in patch *j* when there is no dispersal and no competitors. For slowly dispersing populations, $E_1$ is close to the product ($E_1^1$ ,…,$E_1^m$), and the matrix $A^2$ ($E_1$) is close to the diagonal matrix diag($f^i_2(E_1{}^j)$)$^m_{j=1.}$ Hence the invasion rate $\lambda_2$ ($E_1$) is close to $\max_j f_2^j(E^j_1)>0$. So we obtain: The system with two slowly dispersing competing species is robustly permanent if and only if system refers to system of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

(8) $\max_j f_2^j(E^j_1)>0$ and $\max_j f_2^j(E^j_2)>0$

Intuitively, robust permanence requires for each species there is at least one patch where it can persist.

For *systems, particle interactions, financial transactions, cosmological interactions and reactions, as we shall show model by model the* situation are even easier. Since (under mild assumption on the prey dynamics) the Morse decomposition of $\partial$C is given by two equilibria {0,$E_1$}, the spatial predator-prey system is robustly permanent, if the predator can invade at the prey equilibrium: $\lambda_2$ ($E_1$) > 0. However, the global dynamic are likely to be more complicated. For instance, if in each patch there is a globally stable limit circle (e.g.,[39-41]) and dispersal is sufficiently slow (i.e. $0< e^j{}_i \ll 1$for all *i, j*), our result about robust permanence and the theory of normally hyperbolic manifolds [42] imply that there is a positive *m* dimensional torus, which attracts almost all orbits in C, whenever $\max_j f_2^j(E^j_1)>0$

**5.2 Rock-Paper-Scissors Dynamics.** The Lotka-Volterra model of rock-paper-scissor dynamics is a simple model that is used as property for understanding intransitive ecological outcomes [43,44]. Here, a simple spatial version of this dynamic is given by

$\frac{dx_1^j}{dt} = x_1^j(1 - x_1^j - \beta^j x_2^j - \alpha^j x_3^j) + \sum d_1^{jk} x_1^k - e_1^j x_1^j$

$\frac{dx_2^j}{dt} = x_2^j(1 - \alpha^j x_1^j - x_2^j - \beta^j x_3^j) + \sum d_2^{jk} x_2^k - e_2^j x_2^j$

$\frac{dx_3^j}{dt} = x_i^j(1 - \beta^j x_1^j - \alpha^j x_2^j - \beta^j x_3^j) + \sum d_3^{jk} x_3^k - e_3^j x_3^j$

Where $a^j \in (0,1)$ and $\beta^j > 1$ for all $j$. A more general version of this dynamic is presented in [44]. Under the assumption that there is no cost to dispersal (i.e. $e^j_i = \sum_l d^{l,j}_i$ for all $i, j$) , the maximal invariant set in $\partial C$ consist of the origin 0 and a heteroclinic cycle connecting positive single species equilibria $E_1$, $E_2$ and $E_3$. For slowly dispersing population ($0 < e^j_i \ll 1$ for all $i$, j), $E_1$ is close to $(x_1, x_2, x_3) = (1,\ldots,1,0,\ldots,0,0,\ldots,0)$ and the invasion rates $\lambda_2(E_1)$ and $\lambda_3(E_1)$ are close to $\max_j 1 - a^j$ and $\max_j 1 - \beta^j$. For the other equilibria $E_i$, the debitable rates of the missing species are also given by $\max_j 1 - a^j$ and $\max_j 1 - \beta^j$. Consider the Morse decomposition of $\partial C$ given by 0 and the heteroclinic cycle. Since $\lambda i(0) = 1 > 0$ for $i = 1, 2, 3, 0$ is robustly insaturated. An algebraic computation reveals that (4) holds for all invariant supported by the heteroclinic cycle (i.e. all convex combination of the Dirac measures supported by the equilibria) if and only if

$$\max_j(1 - a^j) > -\max_j(1 - \beta^j)$$

Equivalently,

$$2 > \min_j a^j + \min_j \beta^j$$

In particular, even if the heteroclinic cycle is attracting for each patch when the system is uncoupled (i.e. $2 < a^j + \beta^j$ for all $j$ ), it can be repelling for the weakly coupled system.

5.3. Three species **Lotka Volterra in spatially homogenous environments.** Consider Lotka-Volterra dynamics in a spatially homogenous environment: $(x + a)^n =$

(9) $$\frac{dx_i^j}{dt} = x_i^j\left(r_i - \sum_s a_{is}x_s^j\right) + \sum_l d_i^{jl}x_i^l - e_i^j x_i^j, \qquad i = 1,\dots,k;\ j = 1,\dots,m$$

We assume that dissipative can be shown by a linear Liapunov function: There are $c_i > 0$ such that $\sum_i c_i a_{is}x_i x_s < 0$ holds for all $x \in R_+^k \setminus \{0\}$. Then the weighted sum of densities across all patches $\sum_{i,j} c_i x_i^j$ is decreasing for large densitities and hence (9) generates a dissipative semiflow with global attractor $r$. We further assume that for each patch *j* and each species *i* the immigration rate equals the emigration rate , i. e,

(10) $$e_i^j = \sum_l d_i^{jl}$$

This guarantees that the set H= $\{x \in C: x_i^j = x_i^l \forall i,j,l\}$ of spatially homogenous states is forward invariant under (9).

Now consider the Lotka-Volterra dynamics without spatial structure

(11) $$\frac{dx_i}{dt} = x_i\left(r_i - \sum_s a_{is}x_s\right), \qquad i = 1,\dots,k$$

(11) $$\frac{dx_i}{dt} = x_i(r_i - a_{is}x_s), \quad i = 1,\dots,k$$

It has been shown [12,ch.16.1,16.2] That for *k* =3,(11) is robustly permanent (robust meaning here within the class of Lotka –Volterra systems) if and only if all equilibria on the boundary of $R_+^3$ are unsaturated and whenever there is a heteroclinic cycle connecting the one species equilibria then this cycle is repelling. Equivalently the following four conditions hold. There exists an interior equilibrium $\dot{x}$ (i.e $A\dot{x} = r \ \ with\ \dot{x} \gg 0$); Det (-A)>0; the 2 species subsystems are not bistable competition systems if there is a heteroclinic cycle between the one species equilibria, say $\mathbf{E_1 \to E_2 \to E_3 \to E_1}$ then the following inequality holds

12) $$\lambda_2(E_1)\lambda_3(E_2)\lambda_1(E_3) > |\lambda_3(E_1)\lambda_1(E_2)\lambda_2(E_3)|$$

Furthermore, the proof in [12, ch.16.1, 16.2] shows that the boundary flow of (11) has a simple Morse decomposition: if (iv) applies then there are two Morse sets, the origin 0 is a repeller, and the heteroclinic cycle is the dual attractor (within$\partial R^3_+$); if (iv) does not apply then the finitely many boundary equilibria form a Morse decomposition of the boundary flow.

**Species, population we mean types of and population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.**

**Theorem 4:** Under the above assumptions,(9) for k=3 is permanent if and only if *(i),(ii),(iii)* and *(iv)* hold.

Proof: Suppose (9) is robustly permanent. Since the nonspatial system (11) is the restriction of (9) to the invariant subspace **H**, (11) must be robustly permanent and hence*(i)-(iv)* hold

Conversely, suppose that (i)-(iv) hold. We show that the above Morse decomposition of the boundary flow of (11) (now in **H**) is also Morse decomposition for (9) restricted to the maximal invariant subset of $\partial C$ which is contained in the spatial two species subsystems of (9). For this we show that in each 2 species system the orbits of (9) converge to **H** (For the one species subsystems this is obvious.) If a two species subsystem of (11) has an internal equilibrium, say $\mathbf{E_{12}}$ then by assumption (ii),this is not a saddle but a sink, and hence $a_{11}a_{22} > a_{12}a_{21}$. This implies that the 1-2 submatrix of A is VL-stable [12,ch.15.3,],and [45] implies that $\mathbf{E_{12}}$ is the global attractor for the spatially structured system (9) restricted to the first two species.

If a two species subsystem of (11) has no internal equilibrium, then a one species equilibrium say $\mathbf{E_1}$ is the global attractor. Then Lemma 2 (see below) show the corresponding result for (9).A similar proof applies to the case where the origin **0** is the global attractor of a two species subsystem.

Finally observe that under the above assumptions (9)-(10), the invasion rates $\lambda_i$ at spatially homogenous boundary equilibria are the same for (9) and (11) Since all boundary equilibria are unsaturated for (11), they are for (9) as well. The same applies in case (iv) to the heteroclinic cycle. Note that the condition (12) on the eigenvalues is equivalent to the fact that all invariant measures supported on the heteroclinic cycle are unsaturated, see [16,Ex.4.5]

Species we mean types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

**Lemma 2:** Suppose, in a two species Lotka –Volterra system (11) (k=2) all interior orbits converge to the one species equilibrium $\mathbf{E_1}$ .Then all orbits of the spatial version (9) (k=2) converge to the spatially homogeneous one species equibrium $\mathbf{E_1} \in \mathbf{H}$.

**Proof:** If the system is of predator-prey or mutualistic type then Hasting's [45] result applies again. So let us assume that the local interaction is given by the competition system

(13) $\frac{dx_1}{dt} = r_1 x_1(1 - x_1 - ax_2)$ $\frac{dx_2}{dt} = r_2 x_2(1 - \beta x_1 - x_2)$

With 0<a <1< $beta$.In this case, Kazuo Kishimoto (letter to JH, Oct 1987) has given the following argument. There exists a family of forward invariant rectangles contracting to $\mathbf{E_1}$. This shows that all interior solutions of (9) (k=2) converge to $\mathbf{E_1}$.

Theorem 4 shows that a spatial network of identical patches is permanent, if the within patch dynamics is permanent, and three species Lotka-Volterra dynamics. It is not clear how to extend this result to more than three species. The crucial step in the three species case is that permanence precludes bistable two species subsystems. The spatial version of a bistable two species system allows plenty of stable equilibria outside **H,** see [33].However, permanent four species Lotka Volterra systems may have two species bistable subsystems [12,ch.16.4].So there may be invariant sets on the boundary outside **H** which need to be unsaturated to make (9) permanent.

## 6. MORE APPLICATIONS

6.1. **Disease, particle, product, marketing, financial, econometric dynamics with density-dependent demography**. GAO and Hethcote [46] introduced a model of disease dynamics with density-dependent demography. Here we consider a variation of their model in which a population of size P has a per-capita birth rate b(P) and per –capita death rate d. To ensure that the population persists in the absence of the disease, we assume that there exists K>0 such that b(K)=d. To describe the disease dynamics, let *S* be the number of individuals susceptible to the disease, *I* the Number of infected individuals, and R the number of individuals that have recovered from the disease. We assume that P=S+*I* +R. On the same lines we can think of sick and non performance accounts, and those which are conducting themselves well, and those which are satisfactory and find a parallel formula. The same thing

holds true in case of a black hole eating matter and black hole radiating radiation.. If the disease transmission is asymptotic, β is the contact rate between susceptible and infected individuals, γ is the rate at which individuals recover from the disease, and m the mortality rate due to disease, then the population dynamics is given by

$$\frac{dp}{dt} = b(P)P - dP - mI$$

$$\frac{dI}{dt} = \beta(P - I - R)\frac{I}{\epsilon + P} - (\gamma + d + m)I$$

$$\frac{dR}{dt} = \gamma I - dR$$

Where $\epsilon>0$ is a constant It is useful to introduce a change of coordinates in which $y=\frac{I}{P}$ and $z=\frac{R}{P}$:

$$\frac{dP}{dt} = (b(P)-d-my)P$$

$$\frac{dy}{dt} = \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(P)\ y$$

$$\frac{dz}{dt} = \gamma y - (b(P) - my)z$$

Setting $x_1$=P and $x_2$=(y,z) 'where' denotes transpose yields a structured model where $A_1(x)$=b-d-my and

$$A_2(x) = \begin{pmatrix} \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(p) & 0 \\ \gamma & my-b(p) \end{pmatrix}$$

Since {0,(*K*,0,0)}is a Morse decomposition of the boundary dynamics Theorem 3 and Remark 1 imply that this model is robustly permanent if and only if b(0)>d and $\beta\frac{K}{\epsilon+K} > \gamma + m + d$. In particular, for $\epsilon > 0$ sufficiently small, one requires that the basic reproductive number $\frac{\beta}{r+m+d}$ is greater then one.

Remarkably this same criterion determines robust permanence of model with significantly more complicated boundary dynamics. For example, suppose the focal population is a predator species. If the prey has abundance *N* and exhibits logistic dynamics f(*N*)is the per-capita predator consumption rate of the prey and b(N)is the per-capita reproductive rate of the predator then the dynamics become

$$\frac{dN}{dt} = rN(1 - N/K - f(N)P$$

$$\frac{dP}{dt} = (b(N) - d - my)P$$

$$\frac{dy}{dt} = \left[ \beta(1 - y - z)\frac{P}{\epsilon+P} - r - m(1 - y) - b(N)\ y \right]$$

$$\frac{dz}{dt} = \gamma y - (b(N) - my)z$$

Where r>0 and *K*>0, and *f*(*N*),b(N)satisfy f(0)=b(0)Define $x_1$=*N*, $x_2$=*p*, $x_3$=*(y,z)*′ Let us assume that the disease-free community (i.e.y=z=0)is permanent (i.e.b(K)>d) and let A be the global attractor in the interior of the P-N plane Under these assumption a Morse decomposition of the boundary dynamics is given by $(x_1, x_2, x_3) = (0,0,0)$, $(x_1, x_2, x_3) = (K, 0,0)$and A. Since the equilibria (0,0,0) and (k,0,0) are unsaturated, it remains to characterize $\lambda_3(\mu)$for any invariant measure μ supported by A. Proposition 3 and Remark 1 imply

$$\lambda_2(\mu) = \int_c (bx_1)\text{-}d)d\mu(x) = 0$$

*Hence* $= \int_c bx_1)d\mu(x) = d$ *and*

$$\lambda_3(\mu) = \beta \int \frac{x_2}{\epsilon+x_2} d\mu(x) - r - m - d$$

**Theorem 3**: and Remark 1 implies that this structured model is robustly permanent if $\frac{\beta}{r+m+d} > 1 and\ \epsilon > 0$ is sufficiently small.

6.2. Gene networks. The repressilator is an oscillatory gene network based on three (or more generally an odd number of transcriptional repressors. Such a system has been genetically engineered in E-coli by Elowitz and Leibler [47] Mathematical model go back to [48.49] they all involve concentrations of proteins and mRNs. A modified model with auto-activation was suggested in [50]

$$(14)\ x_i = \beta(y_i - x_i)$$

$$(15)\ y_i = aF(x_{)i} - y_i)$$

With

$$(16)\quad F(x)_i = x_i g\ (x_i, x_i - 1) = \frac{x_i}{1+x_i+p x_{i-1}+Kp\, x_i x_i - 1}$$

Here $x_i$ and $y_i$ are normalized concentrations of proteins and mRN As belonging to gene $i$,p is the strength of repression and k a cooperativity parameter this is a structured system of type(2)

$$(17)\ \begin{pmatrix} x_i \\ y_i \end{pmatrix} dot = \begin{pmatrix} -\beta & \beta \\ ag(x_i,x_i) & -1 \end{pmatrix} \begin{pmatrix} x_i \\ y_i \end{pmatrix}$$

In [50,Thm 3] it was shown that for n=3,a>1 and p>1,system(14)has a heteroclinic cycle connecting the 3 single gene equilibria (similar to the rock –scissors-paper dynamics of section5.2)Moreover

- the system is permanent i.e. the boundary of $R^6$+is repelling, if $\lambda + \mu > 0$
- this heteroclinic cycle is asymptotically stable if $\lambda + \mu < 0$

where $\lambda > 0$ and $\mu < 0$ are the invasion rates of the two missing genes i.e. the leading eignvalues of the 2X2 matrix in (17)evaluated at the appropriate single gene equilibria (for explicit expressions see[50,eqs.(127,128)].)

Theorem 3 implies the permanence condition implies the stronger conclusion of robust permanence A similar result holds for n odd, but the heteroclinic cycle is then actually a heteroclinic network.

Note: For references please see the original paper.

One of the examples for which we apply the model are as follows: **Progress in Particle and Nuclear Physics Volume 88, May 2016, Pages 1–28 Chiral magnetic and vortical effects in high-energy nuclear collisions—A status report D.E. Kharzeeva, b, J. Liaoc, d, , , S.A. Voloshine, G. Wangf**

The interplay of quantum anomalies with magnetic field and vorticity results in a variety of novel non-dissipative transport phenomena in systems with chiral fermions, including the quark–gluon plasma. Among them is the Chiral Magnetic Effect (CME)—the generation of electric current along an external magnetic field induced by chirality imbalance. Because the chirality imbalance is related to the global topology of gauge fields, the CME current is topologically protected and hence non-dissipative even in the presence of strong interactions. As a result, the CME and related quantum phenomena affect the hydrodynamical and transport behavior of strongly coupled quark–gluon plasma, and can be studied in relativistic heavy ion collisions where strong magnetic fields are created by the colliding ions. Evidence for the CME and related phenomena has been reported by the STAR Collaboration at Relativistic Heavy Ion Collider at BNL, and by the ALICE Collaboration at the Large Hadron Collider at CERN. The goal of the present review is to provide an elementary introduction into the physics of anomalous chiral effects, to describe the current status of experimental studies in heavy ion physics, and to outline the future work, both in experiment and theory, needed to eliminate the existing uncertainties in the interpretation of the data. Keywords

Chiral magnetic effect; Chiral vortical effect; Chiral anomaly; Quark-gluon plasma; Heavy ion collisions Corresponding author at: Physics Department and Center for Exploration of Energy and Matter, Indiana University, 727 E Third Street, Bloomington, IN 47405, USA. Copyright © 2016 Elsevier B.V. All rights reserved.

## Sufficient Conditions For The Existence And Linear Stability Of Time Independent Age Distributions In Respect Of Particles, Propositions, Problems, Hegelian Structures, Deleuzean Treatises, Interactions And Transactions In Cosmos And Others

**Abstract:** Gopalaswamy studied sufficient conditions are obtained for the existence and linear stability of time independent age distributions in two species competition with age time lagged density dependent mortality and fertility functions.. Work was published in Bulletin Australian Mathematical Society Volume 25 (1982). We show how the model could be applied in case of finance, econometrics, psycho economics, interactions between particles and transactions in nature. Basically the competition is for resources, notwithstanding such experiments could be conducted in laboratory like a rat kept in front of a cat, such incidents could be represented mathematically, and the corresponding linear stability conditions ca be obtained in case of interactions in cosmos. Take for example the case of a cash credit account of a borrower. Many inward bills waiting for paid, or inland or foreign bills payment all are in line for the debit of the account, which also acts as a resource. In addition there are recurrent cases where in the credit is issued to the cash credit account in case of demand and inland bills. In case of economics, investment has inventories and it has to be utilized towards the end of consummation of output. And this is case of output gobbling part of investment which is inventories. Output has to be sold to the customer and the amount received is diverted partly to savings and consumption and part could be employed towards the consolidation of the system. So this is a case of investment itself acting as a both receiver and payer. It shall eat and also produce the finance. Savings as stated above could be deployed for the further diversification and proliferation and that is towards the end of putting the company on terra firma. Another example in cosmology is the case of black holes guzzling matter. There are lot of black holes and lot of matter. Result of such a guzzling the production of Hawking radiation. Horava gravity is another example. How the conformal field theory could be changed in the case of Horava gravity. Still another case is the case of Vonneumann algebra in the cae of Horava gravity. In this case like cash credit account, the Horava gravity and Von Neumann algebra act as both receivers and takers. Such incidents are legion in cosmology and even in philosophy. We have considered self as individual general ledger, and the debits and credits are raised to the same account thus anointing as both a receiver and payer. We give plethora of such cases and apply the model. Due to spatial constants, same module is used for all the modules and we rewrite to state that the model is applicable to all modules. This is due to spatial restrictions. What is to be done is that each and every parameter is to be changed with a subscript or a superscript and finally it is to be concatenated towards the end of obtention of prediction model or optimizing model, for Duality of motion model for further analysis. Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation

of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

**Introduction**

Let $P_1$ *(t)* and $P_2$ *(t)* denote the total particles, interactions, transactions, etc, which we dilate upon further investigation. Perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. **Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems** sizes (or biomasses) at time $t\geq 0$ of two interacting species living in a common habitat and competing for a common pool of resources. The competition for resources will be implicit in our model similar to that in the two species Lotka- Volterra competition system. Assuming constant sex ratios in the two species we can consider $P_1$ *(t)* and $P_2$ *(t)*to be the Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

population of f; immigration, emigration and internal dispersion in the habitat are assumed to play no significant role in the dynamics of the community.Popoulation we mean Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

We suppose that the two species contain respectively $p_1\,(a, t)\,da$ and $p_2\,(a, t)\,da$ individuals with ages between $a$ and $a + da\ (a\geq 0)$ at time $t$ so that we have

(1.7) $P_1\,(t) = \int_0^\infty \rho_1(a, t)\,da,\ P_2\,(t) = \int_0^\infty \rho_2(a, t)\,da,\ t\geq 0.$

The rates of change of the two species densities are by definition $D\rho_1\,(a, t)$ and $D\rho_2\,(a, t)$ where

(1.8) $D\rho_i\,(a, t) = \limsup_{h\to 0} \frac{\rho_i(a+h,t+h)-\rho_i(a,t)}{h},\ i = 1, 2.$

Assuming the existence of the limits in (1.2) we consider the following time lagged model system

(1.9) $$\begin{cases} D\rho_1\,(a,t) + f_1\,(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_1(a, t) = 0, & a>0,\ t>0 \\ D\rho_2(a,t) + f_2(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_2(a, t) = 0,' \quad (1.4) \\ \rho_1(0, t) = \int_0^\infty b_1\,(a, P_1\,(t), P_2\,(t-\tau))\rho_2(a, t)da\,, & t>0, \\ \rho_2(0, t) = \int_0^\infty b_2\,(a, P_1\,(t), P_2\,(t-\tau))\rho_1(a, t)da\,, \end{cases}$$

(1.5) $\rho_i\,(a, s) = \varphi_i\,(a, s),\ a.0,\ s\in [-\tau, 0],$ $i = 1, 2,$

(1.6) $P_i(s) = \int_0^\infty \varphi_i\,(a, s)\,da,\ s\in [-\tau, 0],$

Where $\tau$some is fixed nonnegative constant; $f_1$ and $f_2$ denote the age and density dependent mortality rates with time lags in interspecific interactions, $b_1$ and $b_2$ denote the age and density dependent fertility functions again having time lags in the interspecific interactions; $\varphi_1$ and $\varphi_2$ denote the initial age distributions needed for the formulation of the model. Note we mean the population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

The model system (1.1)-(1.6) will be meaningful only if $\varphi_{1,}$ $\varphi_2$, $f_1$, $f_2$, $b_1$, $b2$ are nonnegative functions of the respective arguments. Also since the total initial populations have o be finite, $\varphi_1(\cdot, s)$ should belong to $L_1(R^+)$ for each $s\epsilon$ $[-\tau, 0]$. That $f_i$ and $b_i$ $(i = 1, 2)$ are independent of t explicitly and this assumption can be interpreted to represent the temporal uniformity of the environment.

One of the fundamental questions for (1.1)-(1.6) is the following; under what conditions on $f_{i,}$ $b_i$ and $\tau$ does the system (1.1)-(1.4) have time independent nontrivial solutions, and when such solutions, and when such solutions exist, are they stable with respect to some suitable stability criterion? When such stable time invariant solutions exist for (1.1)-(1.4) we say that the two types has a stable age distribution. Existence of a stable age distribution is the analogue of the existence of a stable steady state for the corresponding age independent system. Types we mean types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

In this article we investigate the above question for the case of a competitive interaction with some additional constraints on the vital rates $f_i$ and $b_i$ $(i = 1, 2)$. Age dependent population systems without time delays have been considered by several authors (Gurt in and MacCamy [3, 4, 5], Gurt in and Levine [6], Haimovici [7] and Rotenberg [8]). The question of the relation between an age independent system and a corresponding age dependent system has been considered by Gurt in and MacCamy [4] and Gurt in and Levine [6] who have established an asymptotic (*as* $t\rightarrow\infty$ ) relation between such models by constructing a higher dimensional 1 umped parameter system (in terms of ordinary differential equations) to represent the age dependent distributed parameter system. Assuming the existence of stationary age distributions Haimovici [7] considers their stability in a system of two interacting populations explicitly taking into consideration the dynamical nature of the encironments resources and pollution.

In this article we add another realistic feature namely time lags in the age dependent models and obtain sufficient conditions for the existence of stationary age distributions in (1.1)-(1.4) and show that the time lags in interspecific

interactions have no effect on the linear stability of the age distributions although the decay rates of perturbations will depend on the magnitudes of the time delays. Such "harmless" natures of time lags in interspecific interactions have been noted by the author for the age independent models in [1, 2].

## 3. Existence and uniqueness of solutions

We first narrate our assumptions on the vital rates of $f_i$ and $b_i$ $(i = 1, 2)$ so that the system (1.1)-(1.6) will denote a competition system. Let $BC\ (0, \infty)$ denote the linear space of bounded continuous real functions on $(0, \infty)$ with the norm defined by

$$||f|| = \{max\ |\ f(t)\ |,\ t\ \in (0, \infty)\}.$$

Let $K_1$ be the cone of nonnegative functions of $BC\ [0, \infty]$. Let $C\ [-\tau, 0]$ denote the space of real continuous functions with a norm defined by $||g|| = \{\ max\ |g(s)\ |,\ s \in [-\tau, 0]\}$; let $K_2$ be the nonnegative cone of functions in $C\ [-\tau, 0]$ . We present our first set of assumptions on $f_i$, $b_i$ and $\varphi_i$ for $i = 1, 2$.

$A_1$. $f_1 \in C\ (R^+ \times K_1 \times K_2, R^+)$, $f_1 \geq d_1^*$ on $R^+ \times K_1 \times K_2$

$f_2 \in C\ (R^+ \times K_2 \times K_1, R^+)$, $f_2 \geq d_2^*$ on $R^+ \times K_2 \times K_1$ ,

( $d_1^*$ , $d_2^*$ are positive constants);

$f_1$ and $f_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ ; that is, *for* $i = 1, 2$,

$$|f_i\ (a, P_1^{'}\ , P_2^{'}) - f_i(a, P_1^{"}\ , P_2^{"})| \leq d_i\ [||P_1^{'} - P_1^{"}, P_2^{'} - P_2^{"}||]$$

($d_1$, $d_2$ are positive constants),

$$\partial f_1/\ \partial P_1, \partial f_1/\ \partial P_2 \in C(R^+ \times K_1 \times K_2, R^+)\ ,$$

$$\partial f_2/\ \partial P_1, \partial f_2/\ \partial P_2 \in C(R^+ \times K_2 \times K_1, R^+),$$

$A_2$. $b_1 \in C\ (R^+ K_1 \times\ K_2, R^+)$, $b_1 \leq \beta_1^* on\ R^+ \times K_1 \times K_2$ ,

$b_2 \in C\ (R^+ K_2 \times\ K_1, R^+)$, $b_2 \leq \beta_2^* on\ R^+ \times K_2 \times K_1$ ,

*($\beta_1^*$ , $\beta_2^*$ are positive constants);*

$b_1$ , $b_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ *; that is for* $i = 1, 2$,

$$|b_i\ (a, P_1^{'}, P_2^{'}) - b_i\ (a, P_1^{"}, P_2^{"})| \leq \beta_i\ [||P_1^{'} - P_1^{"}|| + ||P_2^{'} - P_2^{"}]$$

$B_1$ , $B_2$ Are positive constants),

$b_i(\circ,(P_1,P_2)\epsilon L^{\infty}(R^+,R^+)$ For fixed $P_1,P_2,i=1,2,$

$$\frac{\partial b_1}{\partial P_2},\frac{\partial b_1}{\partial P_2}\epsilon\ C\ R^{+}\text{x}\ k_{\ 1}\text{x}\ k_{\ 2},\text{R}^{-}),$$

$$\frac{\partial b_2}{\partial P_1},\frac{\partial b_2}{\partial P_2}\epsilon\ C\ R^{+}\text{x}\ k_{\ 2}\text{x}\ k_{\ 1},\text{R}^{-}),$$

$A_3.\ \varphi_i\ (\cdot\ ,\text{s})\in \text{L}^{\ 2}\ (\text{R}^{+},\text{R}^{+})\ \cap \text{C}(\text{R}^{+},\text{R}^{+})$ for each $s\in\ [-\tau,0],$

$i,\ =1,\ 2,$

By a solution of $(1.1)-(1.6)$ on $[0,\text{T}]$, we mean

$\text{P}_i\text{: R}^{+}\ \square[-\tau,\text{T}]\rightarrow \text{R}^{+}\ \ (i,\ =1,\ 2,)$ With the following properties:

$S_1.$ $\text{P}_i(\cdot\ ,t)\in \text{L}^{\ 1}\ (\text{R}^{+},\text{R}^{+})$ for each $t\ \in[-\tau,\text{T}];$

$S_2.$ $\text{P}_i(a,\cdot\ )\in \text{C}\ ([-\tau,\text{T}],\text{R}^{+})$ for each $a\ \in \text{R}^{+}$;

$S_3.$ $\text{DP}_i$ Exists along the characteristics $t-a=$ constant

On $\text{R}^{+}\ \square\{\text{R}^{+}\cap\ [0,\text{T}]\}$ And is continuous $\ (i,\ =1,\ 2,);$

$S_4.$ $\text{P}_i$ Satisfies the system (1.1)- (1.6) for

$(a,t,)\ \in \text{R}^{+}\ \square[-\tau,\text{T}]$ .

Let us first convert the system (1.1)- (1.6) into an equivalent system of integral equations. For brevity we denote by P the pair ( $\text{P}_1,\text{P}_2$) If we let

$$\text{P}_i(a+x,t+x)=\ \ \text{P}_i(x),(i,\ =1,\ 2,)\ ,$$

Then (1.3) considered at $(1.3)$ considered at $(a+x,\ t+x)$ becomes

$$\frac{dp_{\ i}(\text{x})}{dx}+f_i(a+x,\ P(t+x))\ \ P_i(x)=0$$

Which has a solution of the form?

$$P_i(x)=P_i(0)\ \exp\left[-\int_0^{\infty} f_{\ i}\big(a+s,P(t+S)\big)ds\right]$$

And hence

$$(2.1)\ p_{\ 1}\ (a+x,t+x)=p_{\ 1}\ (a,t)\ exp\big[-\int_o^x\ f_{\ 1}\ \big(a+s,p(t+s)\big)ds\big]$$

For $a \geq t$ *we choose (a,t)=(a-t,0)* in (1.3) and *x=t* in (2.1) so that

(2.2) $p_1(a,t) = \varphi_1(a-t,0)\exp\left[-\int_o^t f_1\left(a-t+s, p_1(s)\right)ds\right]$

$$O \leq t \leq a, i = 1,2,$$

With

$$P_1(s)=(p_1(s),p_2(s-t)),$$

$$P_2(s)=(p_1(s-t),p_2(s)),$$

For $0 \leq a < t$ we choose *(a,t)=(0,t-a)in (1.3)and x=a in (2.1)* so that

*(2.3)* $p_i(a,t)= B_i(t-a)\exp[-\int_o^a f_1(s, p_i(t-a+s))ds]$

$$t > a \geq, i = 1,2,$$

Where

$$(2.4)\ B_i(t) = p_i(0,t), i = 1,2, t > 0.$$

Define $M_i and\ L_i\ (i = 1,2)$ as follows:

$$M_i(a,t,p_i(t)) = exp[-\int_o^{t-a} f_1(s, p_i(a-a+s))ds]$$

$$(2.5)\ L_i\left(a,t,\ p_i(t)\right) = \exp[-\int_o^t f_1(\text{a+s},\ p_i(s))ds]$$

It will now follow from (1.1),(1.3),(2.2)-2.5)that

$$(2.6)\ p_i(t) = \int_o^t B_1(a)\, M_1(a,t,p_i(t)da + \int_o^\infty \varphi_i(a,0)\ L_i(a,t,p_i(t)da,$$

$$t > 0, i = 1,2,$$

$$(2.7)\ p_i(s) = \Phi_i(s) = \int_o^\infty \varphi_i(a,s)da, s \in [-\tau, 0], i = 1,2,$$

(2.8) $\beta_i(\text{t}) = \int_0^t \beta_i(\text{a}) b_i(\text{t-a}, P_i(t)) M_i.\left(a,t,P_i(t)\right)da$

$$+\int_0^\infty b_i\left(a+t, P_i(t)\right)\varphi_i(a,0)L_i\left(a,t,P_i(t)\right)da$$

$$i = 1, 2, \text{t} > 0$$

the equivalence of (2.6)-(2.8) with (1.1)-(1.6) is established by the following whose proof is identical to a similar result of Gurtin and MacCamy [4]; hence we omit the details of the proof of the following.

THEOREM 1 : *Suppose the system* (1.1)-(1.6) *satisfies the conditions* $A_1, A_2, A_3$ . *If* $(P_1, P_2)$ *is a solution of*

*(1.1)-(1.6) then the total population* $P_1, P_2$ *and the birth rates* $B_1, B_2$ *satisfy the integral equation (2.6)-(2.8).*

*conversely if* $P_1, P_2, B_1, B_2$ *are nonnegative continuous solution of (2.6)-(2.8) on [ 0, T] and if* $P_1, P_2$ *are*

*defined by (2.2)-(2.4) on* $R^+$ x [0, T] *then such* $P_1, P_2$ *provide a solution of (1.1)-(1.6) on* $R^+$ x [0, T] .

The following a *priori* estimates are useful to prove our existence theorem below. It will immediately follow from the bounds for $f_i$ and $b_i$ that

$$M_i(a,t,P) \leq \exp[-d_i^*(t-a)] \leq 1, if \quad t \geq a,$$

*(2.9)*

$$L_i(a,t,P) \leq \exp[-d_i^* t]$$

From (2.7)-(2.8) we derive that

$$\beta_i(t) \leq \int_o^t \beta_i^* \beta_i(a) \exp[-d_i^*(t-a)]\, da + \beta_i^* \Phi_i(0) \exp[-d_i^* t], \quad i = 1,2,$$

And hence by Gronwall's inequality,

(2.10) $$\beta_i(t) \leq \beta_i^* \Phi_i(0) \exp[\delta_i t], \ \delta_i = \beta_i^* - d_i^*$$

(2.5), (2.6) and (2.10) lead to

(2.11) $$|P_i(t) - 0_i(0)| \leq \beta(0) exp[\delta_i a - d_i^*(t-a)]\, da$$

$$+\int_0^\infty \varphi_i(a,0)|1 - exp(-d_i^* t)|da$$

$$\leq \beta_i^* \Phi_i(0) exp(\beta_i^* t)\left\{\frac{1-exp[-(d_i^*+\beta_i^*)t]}{d_i^*+\beta_i^*}\right\} + \beta(0) d_i^* t$$

$$\leq [\beta_i^* \beta(0) exp(\beta_i^* t) \beta(0) d_i^*] t \quad for \ t > 0.$$

We can now prove the following result on the existence and uniqueness of solution of the integral equations (2.)-(2.8)

**THEOREM 2:** *Assume that the system* (1.1)-(1.6) *satisfies the hypotheses*

| |
|---|
| $A_1, A_2, A_3$. *Then there exists a unique set of continuous function* $(P_1, P_2, B_1, B_2)$ such *that* |
| $P_i : R^+ \to R^+, \qquad B_i : R^+ \to R^+ \ (i = 1,2\ ),$ |
| *Which satisfy the integral equations* (2.6)-(2.8) |
| Proof. We will first show the existence of local solution on same (possibly small) interval [0, T] $\subset R^+$ such that |
| $P_i \in \mathrm{C}([0,T], R^+), \quad B_i \in \mathrm{C}([0,T], R^+)\ , i = 1, 2,$ |
| Satisfying (2.6)-(2.8) and then show the $P_i$ and $B_i$ can be continued as solution for any finite interval in $R^+$ for some fixed |
| Positive number $T$ we let |
| $C^+[-\tau, T] = \{f \in \mathrm{C}([-\tau, T], R)\,|f \geq 0\}$ |
| Let $\|.\|_T$ denote the supremum norm in $\mathrm{C}^+[-\tau, T]$. Then for any |
| $(x, y) \in \mathrm{C}^+[-\tau, T] \times \mathrm{C}^+[-\tau, T]$ we define |
| $\|(x,y)\| = \|x\|_T + \|y\|_T,$ |
| $\mathrm{C}_T = \mathrm{C}^+[-\tau, T], \quad \mathrm{C}_T^2 = \mathrm{C}^+[-\tau, T] \times \mathrm{C}^+[-\tau, T].$ |
| |
| Members of $\mathrm{C}_T$ will be denoted by $Z_t$ with the meaning that |
| $Z(t+s) = z_t(s)$ *for* $s \in [-\tau, T]$ *and* $t \in [0, T]$. Now for any fixed |
| $P_t = (P_{1t}, P_{2t}) \in \mathrm{C}_T^2$, (2.8) is a linear system of uncoupled Volterra integral equations in B = $(B_1, B_2)$ and hence (2.8) will |
| have a unique continuous solution under the assumption $A_1 - A_3$. let us denote such a solution by |
| $\mathrm{B}(t) = B^T(P_t), \quad t \in [0, T]$ |
| Or |
| $B_1(t) = B_1^T(P_{1t}, P_{2t}), \quad B_2(t) = B_1^T(P_{1t}, P_{2t}), \ t \in [0, T]$ |
| If we supply this solution $(B_1, B_2)$ in (2.6) we find that (2.) will be satisfied if and only if $(P_{1t}, P_{2t})$ is a fixed point of the |
| operator |

$$\Pi(P_t) = \{\Pi_1(P_{1t}, P_{2t}),\ \Pi_2(P_{1t}, P_{2t})\}$$

Where $\pi_1$ *and* $\pi_2$ respectively denote the right sides of (2.6) for *i =1,2* .

It is not difficult to see from the nature of $M_i$ , $L_i$ and the hypotheses on $f_i, b_i, \varphi_i$ that

$$B: C_T^2 \to C_T^2 \text{ and } \Pi C_T^2 \to C_T^2$$

(which guarantees the nonnegativity of $P_i$ and $B_i$). We will first show that

There exist positive constants $r^*, t^*$ such that $\pi$ is a mapping of a sphere

$$S_{t*}(\Phi, r^*) = \left\{ (P_{1t}, P_{2t}) \in C_{t*}^2 \mid \|P_t - \Phi_t\|_{t*} \le r^* \right\}$$

Into itself and also is a contraction where $\Phi_t$ is defined on [-$\tau$, $t^*$]

By the following:

$$\Phi_t = (\Phi_{1t}, \Phi_{2t});$$

(2.12)

$$\Phi_{it} = \begin{cases} \int_0^\infty \varphi_i(a, s)\, da \text{ for } s \in [-\tau, 0]. \\ \int_0^\infty \varphi_i(a, 0)\, da \text{ for } s \in [0, t^*]. \end{cases}$$

We note that similar to (2.11) one gets the estimates

$$|\Pi(P)(t) - \Phi_i(0)| \le \{B_i^* \Phi_i(0) \exp[B_i^* t] + \Phi_i(0) d_i^*\} t$$

Which together with (2.7) and (2.12) lead to

(2.13) $$\|\Pi(P_t) - \Phi_t\|_{t*} \le \left( \sum_{j=1}^2 \|\Phi_{jt}\|_{t*} [B_j^* \exp(B_j^* t) + d_i^*] \quad t^* \right)$$

Thus if we choose,

$$r^* = T \sum_{j=1}^2 \|\Phi_{jt}\|_{t*} [B_j^* \exp(B_j^* T) + d_i^*]$$

Then $\Pi$ is a mapping of $S_{t*}(\Phi, r^*)$ into itself for all $t^* \le T$. we then will have to show that for small $t^* > 0$ , $\Pi$ is a contraction.

Let $P_t^{(1)}$ *and* $P_t^{(2)}$ be arbitrarily chosen in $S_{t*}(\Phi, r^*)$ and consider $|\Pi(P_t^{(1)}) - \Pi(P_t^{(2)})|$ for $t \le t^*$.

Corresponding to the chosen $P_t^{(1)}$ , $P_t^{(2)}$ we let $B^{(1)}(a) = B(P_T^{(1)})$ and $B^{(2)}(a) = B(P_T^{(2)})$ . Then from (2.6),

$$|\Phi(P^{(1)}(t)) - \Phi(P^{(2)}(t)| \leq \left| \int_0^t B_1^{(1)}(a)\left\{M_1\left(a,t,P_t^{(1)}\right) - M_1\left(a,t,P_t^{(2)}\right)\right\}da \right|$$

$$+ \left| \int_0^t [B_1^{(1)}(a) - B_2^{(2)}(a)]\, M_1(a,t,P_t^{(2)})da \right|$$

$$+ \left| \int_0^t \varphi_1(a,0)\left[L_1\left(a,t,P_t^{(1)}\right) - L_1\left(a,t,P_t^{(2)}\right)\right]da \right|$$

$$= J_1 + J_2 + J_3 \text{ (say)}$$

using the elementary inequality

$$|e^x - e^y| \leq |x-y| \ for \ x,y \leq 0$$

We derive from (2.5) that

$$\left| M_1(a,t,\ p_t^{(1)}) - M_1(a,t,\ p_t^{(2)}) \right|$$

$$\leq \int_o^{t-a} \left| f_1(s,\ p_t^{(1)}(a+s)) - f_1(s,\ p_t^{(2)}(a+s)) \right| ds$$

$$\leq d_1 \| p_t^{(1)} - p_t^{(2)} \|_{t^*} \text{(using } A_1\text{)}$$

*Now,using (2.10),*

$$|J_1| \leq \int_o^a \beta_1^* \Phi_1(0)\, e^{\delta 1^a} \left(\int_o^{t-a} d_1 \| P_t^{(1)} - P_t^{(2)} \|_t ds\right)da$$

$$\leq d_1 \beta_1^* \Phi_1(0)\, t^2 \| P_t^{(1)} - P_t^{(2)} \|_t e^{\delta 1^t/2}$$

In a similar way

$$|J_3| \leq \int_o^\infty \varphi(a,o)[\int_o^t \left| f_1(a+s, P_t^{(1)}(s) - f_1(a+s, P_t^{(2)}(s)) \right| ds]da$$

$$\leq d_1 \Phi_1(0) \| P_t^{(1)} - P_t^{(2)} \|_t$$

Now we have to estimate $\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right|$ in order to estimate $J_2$, from (2.8),

$$B_1^{(1)}(a) - B_1^{(2)}(a)$$

$$= \int_o^a \left[B_1^{(1)}(\sigma) - B_1^{(2)}(\sigma) b_1\left(a-\sigma, P_t^{(1)}(a)\right) M_1(\sigma, a, P_t^{(1)}(a) d\sigma + \int_o^a \left[B_1^{(2)}(\sigma)\right.\right.$$

$$\mathrm{X}\left[b_1(a-\sigma, P_t^{(1)}(a) M_1(\sigma, a, P_t^{(1)}(a) - b_1\left(a-\sigma, P_t^{(1)}(a) M_1\left(\sigma, a, P_t^{(2)}(a)\right)\right]\right] d\sigma$$

$$+\int_o^\infty \varphi(\sigma, o)\left[b_1\left(\sigma + a, P_t^{(1)}(a)\right) - b_1\left(\sigma + a P_t^{(2)}(a)\right)\right], L_1\left(\sigma, a, P_t^{(1)}(a)\right) d\sigma$$

$$+\int_o^\infty \varphi(\sigma, o) b_1\left(\sigma + a, P_t^{(2)}(a)\right)\left[L_1\left(\sigma, a, P_t^{(1)}(a)\right) - L_1\left(\sigma, a, P_t^{(2)}(a)\right)\right] d\sigma$$

$$= R_1 + R_2 + R_3 + R_4 \text{ (say)}.$$

By the uniform Lipschitz continuity of $b_1$ in $A_2$ we get

$$\left| b_1\left(a+\sigma, P_t^{(1)}(a)\right) - b_1\left(\sigma + a, P_t^{(2)}(a)\right) \right| \leq B_1 \, /\!/ \, P_t^{(1)} - P_t^{(2)} /\!/ \text{ for } a \, \varepsilon [o, t^*]$$

And hence we have

$$|R_3(a) + R_4(a)| \leq \beta_1 \Phi_1(0) || P_t^{(1)} - P_t^{(2)} ||_{t^*} + \Phi_1(0) \beta_1 d_1 || P_t^{(1)} - P_t^{(2)} ||_{t^*}^{*a},$$

$$|R_2(a)| = [\int_0^a [B_1^{(2)}(\sigma)\{b_1(a-\sigma, P_t^{(1)}(a)) - b_1(a-\sigma, P_t^{(2)}(a))\} M_1(\sigma, a, P_t^{(1)}(a))$$

$$+ b_1(a, -\sigma, P_t^{(2)}(a))\{M_1(\sigma, a, P_t^{(1)}(a)) - M_1(a-\sigma, P_t^{(2)}(a))\}] d\sigma]$$

$$+ b_1\left(a - 0, P_t^{(2)}(a)\right)\{M_1(\sigma, a, P_t^{(1)}(a)) - M_1(a-\sigma, P_t^{(2)}(a))\}] d\sigma]$$

$$\leq \int_0^a [\beta_1^* \Phi_1(0) \, \mathrm{e}^{\delta 1^a} \beta_1 \| P_t^{(1)} - P_t^{(2)} /\!/_{t^*} + \beta_1 \mathrm{d}_1 \| P_t^{(1)} - P_t^{(2)} /\!/_{t^*} (a-o)] do$$

$$\leq \beta_1^* \Phi_1(0) \, \mathrm{e}^{\delta 1^a} \| P_t^{(1)} - P_t^{(2)} /\!/_{t^*} \{\beta_1 a + d_1 a^2/2\}$$

Thus

$$\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right|$$

$$\leq \int_o^a \left| B_1^{(1)}(\sigma)\text{-}B_1^{(2)}(\sigma) \right| B_1^{*d}\sigma$$

$$+ \Phi_1(0) /\!/ P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} \left\{ B_1 + B_2\, d_1\, a + B_1^{*} \left( B_2 a + d_1 a^2/2 \right) e^{\delta 1^a} \right\}$$

$$\leq \beta_1^* \int_o^a \left| B_1^{(1)}(\sigma)\text{-}B_1^{(2)}(\sigma) \right| \text{do} + \left( r^* + /\!/ \Phi_t /\!/_{t*} \right) \| P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} g_1^{(a)}$$

Where $g_1(a) \rightarrow \beta_1$ as $a \rightarrow 0$. By Gronwall's inequality it will follow that

$$\left| B_1^{(1)}(a)\text{-}B_1^{(2)}(a) \right| \leq \left( r^* + /\!/ \Phi_t /\!/_{t*} \right) /\!/ P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} \{ g_1(a) + 0(a) \}$$

Using this estimate for $R_1(a)$ we get

$$|R_2(a)| \leq \left| \int_0^a [B_1^{(2)}(\sigma)\, B_1^{(2)}(\sigma),]\, b_1(a)) - b_1(a\text{-}\sigma, P_t^{(1)}(a))\, M_1(\sigma, a, P_t^{(1)}(a)) do \right|$$

$$\left| B_1^{*} \left( r^* + /\!/ \Phi_t /\!/_{t*} \right) /\!/ P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} 0(a) as\ a \rightarrow 0 \right.$$

*Thus*

$$\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right| \leq (r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ P_t^{(1)} - P_t^{(2)} /\!/_t o(1) as \quad a \rightarrow o$$

And hence

$$\left| \pi_1(P_t^{(1)} - \pi_1\left(P_t^{(2)}\right) \right| \leq ((r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ P_t^{(1)} - P_t^{(2)} /\!/_t o(t) as \quad t \rightarrow o$$

Which together with a similar bound for $\left| \pi_2(P_t^{(1)} - \pi_2\left(P_t^{(2)}\right) \right|$ leads to

$$\| \pi\, (P_t^{(1)} - \pi\, \left(P_t^{(2)}\right) /\!/ \leq (r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ P_t^{(1)} - P_t^{(2)} /\!/_{t*} o(t^*)$$

As $t^* \rightarrow o$ showing that if $t^*$ is small enough, the mapping $\pi$ on

$s_{t^*}(\Phi, r^*)$ is a contraction. Hence the system (2.6)-(2,8) has a unique solution for $o \leq t \leq t^* for\ t^*$ sufficiently small and positive

Now if [0,T] is any finite interval on $R^+$, we can consider [0,T] as a finite union of intervals of length less than or equal to $t^*$ Since the estimate (2.11)is valid for any finite interval we can extend the solution of (2,6)-(2.8) form $[0, t^*]$ to [0,T] and this completes the proof,

3. Stationary age distributions and their linear stability we will now establish the existence of time independent solution

$P_t^*, P_t^*$ Of (1.1)-(1.6): such solutions satisfy

(3.1) $\frac{dp\,P_2^*(a)}{da}+f_i(\mathrm{a}, P_1^*, P_2^*)\,P_i^*(a) = 0,\ \ i = 1,2, a > 0,$

(3.2) $P_i^* = \int_o^\infty P_i^*(a)da, i = 1,\ 2,$

(3,3) $P_i^*=\int_o^\infty b_i(a,\ P_1^*,\ P_2^*)P_i^*(a)da,\ i = 1.\ \ 2$

All solutions of (3.1)are of the from

(3.4) $P_i^*(a) = P_i^*(o)\exp\left[-\int_o^\infty f_i(s, P_1^*, P_2^*)ds\right], i = 1,\quad 2, a \geq 0,$

Which together with (3.3) lead to

$1=\int_o^\infty b_1(a, P_1^*, P_2^*)\exp\left[-\int_o^a f_1(s, P_1^*, P_2^*)ds\right]da,$

(3.5) $1=\int_o^\infty b_2(a, P_1^*, P_2^*)\exp\left[-\int_o^a f_2(s, P_1^*, P_2^*)ds\right]da$

Thus the existence of stationary solutions of (1.1)-(1.6) reduces to the existence of a pair of positive constants $P_1^*, P_2^*$ satisfying (3.5). We can now prove the following,

THEOREM 3. Assume that $f_1$ and $b_i$ satisfy ( $A_1$ ) and ( $A_2$ )and let

(3.6) $F_{i.}\ (P_1, P_2) = \int_o^\infty b_i(a, P_1, P_2)\exp\left[-\int_o^\infty f_i(s, P_1, P_2)ds\right]da,$

$i$ =1, 2, $P_1 \geq 0$, $P_2 \geq 0$,

**And suppose the following hold:**

(3.7) $A_4$. $F_{i.}\ (0,\ 0) > 1,\ i = 1,2$:

$A_5$. *there exists a positive constant c for which*

$F_1$(C,0)< $F_{2.}$(C,0),

(3.8) $F_1$(0.C)< $F_{2.}$(0,C),

Then there exists a unique pair $(P_1^*, P_2^*)$ of real numbers $P_1^* > o$,

$P_2^* > 0$ such that

(3,9) $F_1(P_1^*, P_2^*) = 1 = F_2(P_1^*, P_2^*)$

And the unique nonegative solution of (3.1)-(3.3) is given by

(3.10) $P_2^*(a) = [P_2^* \exp[-f_i(s, P_1^*, P_2^*)ds])/(\int_o^\infty \exp[-\int_o^a f_i(s, P_1^*, P_2^*)ds]da)$,

$i = 1, 2, \quad a \geq 0,$

Proof Consider the elements of the surfaces defined by

$$Z = F_1(P_1, P_2), z = F_2(P_1, P_2), \ P_1 \geq 0, P_2 \geq 0,$$

In the ($P_1$,$P_2$,z) space. By hypothesis ($A_1$) we have $\partial F_i / \partial P_j < 0$ *and* , by ($A_2$), $\partial b_i / \partial P_j \leq 0$ for all

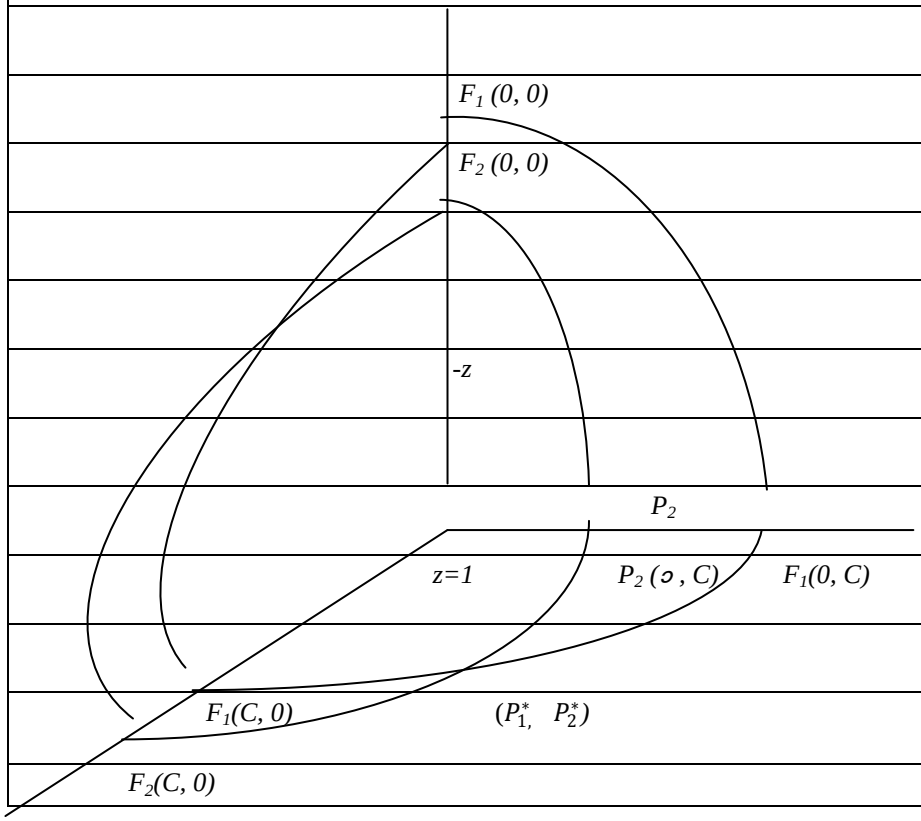
$$P_1 \geq 0, P_2 \geq 0, \quad i = 1, 2,$$

$j = 1, 2,$ Hence by ($A_4$) the intersections of these surfaces with the plane z=1 define two curves on the z=1 plane connecting the I ines {z=1, $P_1 \geq 0, P_2 \geq 0,$} and {z= $1, P_2 \geq 0, P_1 \geq 0$, These two curves are defined by

$$F_1(P_1, P_2 = 1 \ and \ F_2(P_1, P_2) = 1, \ F_1(P_1, P_2 \geq 0,$$

Which by ($A_s$) will intersect at a unique point say $((P_1^*, P_2^*)$ with $(P_1^* > 0, P_2^* > 0$ (see figure). The uniquess of the point $((P_1^*, P_2^*)$ *is a* consequence of the smoothness of the surfaces $z = F_1(P_1, P_2)$ and $z = F_2(P_1, P_2)$

Following is a hypothetical case for Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates attributable or ascribable in the transactions and reactionary potentialities.



(The figure corresponds to the case $F_1(0,0) > F_2(0,0) > 0$. The other possibility $1 < F_1(0,0)$ is treated similarly) It will now

Follow that for such a pair $(p^*_1, p^*_2)$ the unique nonnegative solution of (3.1) - (3.3) is given by (3.10).

Now to examine the linear stability of the stationary age distributions $p^*_1, p^*_2$, we let

(3.11)
$$\begin{cases} p_i(a,t) = p^*_i(a) + u_i(a,t) \\ p_i(t) = p^*_i + p_i(t), \quad i = 1, 2, \quad a > 0, \quad t > 0, \\ p_i(t) = \int_0^\infty u_i(a,t)da, \end{cases}$$

In (1.1)-(1.5) and derive the following variational system after neglecting the nonlinear terms in the perturbations $\boldsymbol{u}_i$ and $\boldsymbol{p}_i$ $(i = 1, 2,)$ (such a procedure of linear stability analysis can be justified As has been done in Gurtin and MacCamy [4].)

(3.12)
$$\begin{cases} Du_i(a,t) = -f_i\left(a, p^*_1, p^*_2\right) u_i(a,t) - p^*_i(a)\left\{\sum_{i=1}^{2} \frac{\partial f_i}{\partial p^*_i} p_i(t)\right\}, \\ u_i(0,t) = \int_0^\infty b_i\left(a, p^*_1, p^*_2\right) u_i(a,t)da + \int_0^\infty p^*_i(a)\left\{\sum_{i=1}^{2} \frac{\partial b_i}{\partial p^*_i} p_i(t)\right\} da \end{cases}$$

Where $i = 1,2$ and (3.13)
$$p_{11}(t) = p_1(t), \quad p_{12}(t) = p_2(t-\tau), \quad p_{21}(t) = p_1(t-\tau), \quad p_{22}(t) = p_2(t)$$

To consider the asymptotic bahaviour as $t \to \infty$ as of solutions of (3.12) we let

(3.14)
$$\begin{cases} u_i(a,t) = u_i(a)\exp[\lambda t], \quad u_i(a) \to 0 \text{ as } a \to \infty, \\ p_i(t) = p^*_i \exp[\lambda t], \\ \qquad\qquad i = 1, 2, \quad a > 0, \quad t > 0, \\ p^*_i = \int_0^\infty u_i(a)da, \end{cases}$$

Using (3.14) in (3.12) - (3.13),

$$\frac{du_i}{da} + \lambda u_i = -f_i\left(a, p^*_1, p^*_2\right) u_i(a) - p^*_i(a)\left\{\sum_{j=1}^{2} \frac{\partial f_i}{\partial p^*_j} p^*_i\right\},$$

(3.15)

$$\mu_i(0) = \int_0^\infty b_i\left(a, p_1^*, p_2^*\right)\mu_i(a)da + \int_0^\infty p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial b_i}{\partial p_j^*}\ p_i^*\right\}da,$$

Where $i = 1,\ 2,$ and

$$p_{11}^* = p_1^*, \quad p_{12}^* = p_2^* \exp\{-\lambda\tau\},$$

$$p_{21}^* = p_1^*,\ \exp\{-\lambda\tau\},\ p_{22}^* = p_2^*.$$

If we integrate the first of (3.15) and use the fact that $\mu_i(a) \to 0$ as $a \to \infty$ we can derive that $p_1^*$ and $p_2^*$ are governed by

(3.16) $$\lambda p_1^* - \int_0^\infty p_1^*(a)\left[\left(\frac{\partial b_1}{\partial p_j^*} - \frac{\partial f_1}{\partial p_1^*}\right)p_1^* + \left(\frac{\partial b_1}{\partial p_2^*} - \frac{\partial f_1}{\partial p_2^*}\right)p_2^* e^{-\lambda\tau}\right]da$$

$$= \int_0^\infty \left[b_1\left(a, p_1^*, p_2^*\right) - f_1\left(a, p_1^*, p_2^*\right)\right]\mu_i(a)da,$$

(3.17) $$\lambda p_2^* - \int_0^\infty p_2(a)\left[\left(\frac{\partial b_2}{\partial p_1^*} - \frac{\partial f_2}{\partial p_1^*}\right)p_1^* e^{-\lambda\tau} + \left(\frac{\partial b_2}{\partial p_2^*} - \frac{\partial f_2}{\partial p_2^*}\right)p_2^*\right]da$$

$$= \int_0^\infty \left[b_2\left(a, p_1^*, p_2^*\right) - f_2\left(a, p_1^*, p_2^*\right)\right]\mu_2(a)da,$$

From the first of (3.15) we get

(3.18) $$\mu_i(a) = \left\{\mu_i(0)p_i^*(a)/p_i^*(0)\exp[\lambda a]\right\} - p_i^*(a)\exp[-\lambda a]\int_0^a \left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right\}e^{\lambda s}ds, \quad i = 1, 2,$$

Which by the second of (3.15) becomes,

(3.19) $$\mu_i(a)\left[p_i^*(a)/p_i^*(0)c_i\exp(\lambda a)\right]\left[\int_0^a p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right\}da\right.$$

$$\left. - \int_0^\infty b_i\left(a, p_1^*, p_2^*\right)p_i^*(a)\,e^{-\lambda a}\left\{\int_0^a\left(\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right)e^{\lambda s}ds\right\}da\right]$$

$$- p_i^*(a)\,e^{-\lambda a}\int_0^a\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right\}e^{\lambda s}ds, i = 1, 2,$$

where

(3.20) $$c_i = 11 - \frac{1}{p_i^*(0)} \int_0^\infty b_i\left(a, p_1^*, p_2^*\right) p_i^*(a) e^{-\lambda a} da, \quad i = 1, 2,$$

Now supplying $\varphi_1(a)$ and $\varphi_2(a)$ from (3.19) in (3.16)-(3.17) we derive that

$$\left(\lambda - A_{11}\right) p_1^* - A_{12} p_2^* \exp\left(-\lambda\tau\right) = g_{11}(\lambda) p_1^* + g_{12}(\lambda) p_2^* \exp(-\lambda\tau),$$

(3.21)

$$-A_{21} p_1^* \exp\left(-\lambda\tau\right) + \left(\lambda - A_{22}\right) p_2^* = g_{21}(\lambda) p_1^* \exp\left(-\lambda\tau\right) + g_{22}(\lambda) p_2^*$$

Where

(3.22) $$A_{i\varphi} = \int_0^\infty p_i^*(a) \left[\frac{\partial b_i}{\partial p_\varphi^*} - \frac{\partial f_i}{\partial p_\varphi^*}\right] da,$$

(3.23) $$g_{i\varphi}(\lambda) = \left[p_i^*(0) c_i\right]^{-1} \left\{\int_0^\infty \sigma_i(a) p_i^*(a) e^{-\lambda a} da\right\} \varphi$$
$$\left\{\int_0^\infty \left[p_i^*(a) \frac{\partial b_i}{\partial p_\varphi^*} - b_i\left(a, p_1^*, p_2^*\right) p_i^*(a) e^{-\lambda a} \left(\int_0^a \frac{\partial f_i}{\partial p_\varphi^*} e^{\lambda s} ds\right)\right] da\right\}$$

$$- \int_0^\infty \sigma_i(a) p_i^*(a) e^{-\lambda a} \left(\int_0^a \frac{\partial f_i}{\partial p_\varphi^*} e^{\lambda s} ds\right) da, \quad i, \varphi = 1, 2,$$

(3.24) $$\sigma_i(a) = b_i\left(a, p_1^*, p_2^*\right) - f_i\left(a, p_1^*, p_2^*\right), \quad i, = 1, 2.$$

To solve the linear perturbational system (3.12) it is enough to find the constants $\lambda, p_1^*, p_2^*$ satisfying (3.21) since one can then use (3.19) to find the perturbations from (3.14). A nontrivial solution $\left(p_1^*, p_2^*\right)$ for (3.21) will exist if and only if $\lambda$ is a root of the equation $et \begin{bmatrix} \lambda - \left(A_{11} + g_{11}(\lambda)\right) & -\left(A_{12} + g_{12}\right) e^{-\lambda\tau} \\ -\left(A_{21} + g_{21}\right) e^{-\lambda\tau} & \lambda - \left(A_{22} + g_{22}(\lambda)\right) \end{bmatrix} = 0$

Or equivalently

(3.25) $$\lambda^2 - \lambda\left(A_{11} + A_{22}\right) + A_{11}A_{22} - A_{12}A_{21} \exp\left[-2\lambda\tau\right] = \lambda s_1(\lambda) + s_2(\lambda, \tau)$$

Where $S_1(\lambda) = g_{11}(\lambda) + g_{22}(\lambda)$ ,

$S_2(\lambda, \tau) = \left[A_{21} g_{12}(\lambda) + A_{12} g_{21}(\lambda) + g_{12}(\lambda) g_{21}(\lambda)\right] \exp\left[-2\lambda\tau\right]$

$$-A_{11} g_{22}(\lambda) - A_{22} g_{11}(\lambda) - g_{11}(\lambda) g_{22}(\lambda)$$

The following result establishes the linear local asymptotic Stability of the stationary age distributions $p_1^*(a)$ and $p_2^*(a)$

THEOREM 4 Assume the following:

$A_6.\ 0 < \beta_i^* < d_i^*\ (i,\ =1,\ 2.)\ \ (see(A_1)\ and\ (A_2));$

$A_7.\ \{A_{11} + A_{22} + \| S_1(0) \|\}^2 < 4\{(A_{11}A_{22} - A_{12}A_{21}) - \| S_2(0) \|\},$

$$\| S_1(0) \| = |g_{11}(0)| + |g_{22}(0)|,$$

$$\| S_2(0) \| = |A_{11}\, g_{12}(0)| + |A_{12}\, g_{21}(0)| + |g_{12}(0)\, g_{21}(0)|$$

$$+ |A_{11}\, g_{22}(0)| + |A_{22}\, g_{11}(0)| + |g_{11}(0)\, g_{22}(0)|$$

Then all the roots of (3.25) have negative real parts and hence the

perturbations $u_i\ (a,\ t) = \phi(a) \exp[\lambda t] \to 0\ as\ \ t \to \infty$ for $i = 1, 2\ and\ a \geq 0$

Proof . Define $F_1(\lambda, \tau)$ and $F_2(\lambda, \tau)$ as follows:

$F_1(\lambda, \tau) = \lambda^2 - \lambda(A_{11} + A_{22}) + A_{11}A_{22} - A_{12}A_{21} \exp[-2\lambda\tau];$

$F_2(\lambda, \tau) = \lambda S_1(\lambda) + S_2(\lambda, \tau),$

Using $(A_6)$ one can show that for $\lambda$ with re $(\lambda) \geq 0$, $c_1$ and $c_2$ are positive and bounded away from zero, Now since $A_{ij} > 0$ and $\tau \geq 0$ we find from the nature of the dependence of $g_{ij}$ on $\lambda$ that for all $\lambda\ \ \mu \pm iw$ with $\mu \geq 0$ and $w \geq 0$ we have

(3.26) $$\left| F_1(\lambda, \tau) - P_2(\lambda, \tau) \right|$$

$$\geq \left| F_1(\lambda, 0) | - F_2(\lambda, \tau) \right|$$

$$\geq |\lambda|^2 - |\lambda|\{A_{11} + A_{22} + \| S_1(0) \|\} + \{A_{11}A_{22} - A_{12}A_{21} - \| S_2(0, \tau) \|\}$$

$$\geq \left[(|\lambda| - P)^2 + (q)\right]^2$$

Where

$p = \{A_{11} + A_{22} + \| S_1(0) \|\}/2,$

$q = \left\{4 | (A_{11}A_{22} - A_{12}A_{21}) - \| S_2(0, T) \| | - | A_{11} + A_{22} + \| S_1(0) \| |^2\right\}/2,$

It will now follow, from (A) and (3.26), $F_1\ (\lambda, \tau) = F_2(\lambda, \tau)$ connot have roots with zero or positive real parts and this completes the proof.

In conclusion we remark that the condition (3.7) means that the net creation rates of each system is greater than utility; a condition of this type is well known in age wise classification of particles, atoms, molecules, cosmological phenomenon. (3.8) can be interpreted to mean that the intraspecific inhibition is higher than that of the

interspecific interaction. The conditions of Theorem 4 are analytical and do not lend themselves for any worthwhile interpretation in terms of the model parameters. So also an interpretation can be recorded as a remark in respect of destruction rates of particles, properties, elements, or projections.

Note: The model is applicable to each and every module. As stated in the foregoing, there are possibilities of situation where an account itself be a receiver and it is also the one to which debit is raises. Such examples are legion in all the fields of finance, economics, political science and physics and chemistry, and of course mathematics. Classification based on time is the preeminently and predominantly suited premise for such model. We repeat that the model is to be read with each module We will expatiate the importance and implicational ramifications as we proceed with the work. Assortment-blend, collection-combinations, combo-gallimaufry, goulash-hash, hodgepodge-medley, mishmash-mixture, motley-pastiche patchwork-salmagundi, soup-stew, mixed bag-mélange, olio- olio podrida, assortment variety, Array choice, collection combination, combo diversity, hodgepodge jumble, kind medley, miscellany mishmash, mixed bag mixture, mélange potpourri, selection sort, brew concoction, beverage blend, broth compound, distillation drink, fermentation hash, hodgepodge infusion, instillation liquor, mélange miscellany, mishmash mixture, mélange potpourri, borscht bouillon, bowl brew, chowder concoction, consommé decoction, gumbo hodge-podge, olio porridge, potage potpourri, pottage puree, splash stock, vichyssoise hodgepodge, mixture- mess, collection combination, goulash hash, jumble medley, miscellany mishmash, mixed bag mélange, olio patchwork, potpourri salmagundi, jumble clutter, disarrangement disarray, disorder farrago, gallimaufry garbage, goulash hash, litter medley, mess miscellany, mishmash mixture, muddle mélange, olio pastiche, patchwork potpourri, salmagundi scramble, shuffle snarl, tangle tumble, medley miscellany, assortment brew, collection combo, conglomeration farrago, hodgepodge jumble, melee mingling, mishmash mixed bag, mixture mélange, pasticcio pastiche, patchwork potpourri, salmagundi variety, mélange mixture, assortment combo, farrago gallimaufry, pasticcio pastiche, patchwork potpourri, salmagundi soup, stew miscellany, accumulation aggregation, anthology assortment, brew collectanea, combination combo, compilation conglomeration, accumulation diversity, farrago gallimaufry, melee mishmash, mix mixed bag, mixture muddle, mélange odds and ends, olio pasticcio, pastiche patchwork, potpourri salad, salmagundi smorgasbord, stew variety, mishmash jumble, medley mess, miscellany mixed bag, mixture mélange, olio patchwork, potpourri salmagundi, mixed bag array, assortment choice, collection combination, combo conglomerate, diversity hodgepodge, jumble medley, mélange miscellany, mishmash mixture, montage mélange, odds and ends, patchwork potpourri, smorgasbord sundries, mix/mixture assortment, combination, admixture adulteration, alloy amalgam, amalgamation assimilation, association batter, blend brew, combine combo, commixture composite, compound concoction, confection conglomeration, cross crossing dough fusion, goulash grab bag, hodgepodge hybrid, hybridization incorporation, infiltration interfusion, jumble mash, medley merger, mingling miscellany, mishmash mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, odds and ends assortment, bits and pieces, debris etcetera, hodgepodge jumble, knickknacks leavings, leftovers litter, medley melee, miscellaneous paraphernalia, miscellany mishmash, mixed bag motley, mélange notions, novelties oddments, olio particles, patchwork potpourri, rags refuse, remainder remains, remnants rest, rummage scraps, sundries sundry items, compilation mishmash, paste-up patchwork, potpourri reappropriation, reproduction synthesis, conglomeration cross, crossing dough, fusion goulash, interfusion jumble, mash medley, merger mingling, miscellany mishmash, mix mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, colluviums gallimaufry, hash mixture, pasticcio pastiche, conglomeration cross section, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multiformity multiplicity, mélange potpourri, range shift, variegation variousness, divergencies diverseness, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multifariousness multiformity,

mélange potpourri, range shift, soup stew, variance variation, variegation variousness, oddment odds and ends, assortment bits, polymorphism assortment, change collection, combo conglomeration, cross section departure, discrepancy disparateness, divergencies diverseness, diversification diversity, fluctuation heterogeneity, heterogeneousness incongruity, intermixture many-sidedness, medley miscellaneousness, miscellany mishmash, mixed bag mixture, notions novelties, oddment oddments, olio particles, patchwork potpourri of the particles, interactions, accounts, transactions in the cosmos. The wisdom and understanding of our own divine essence is like a beaconing begorra, bureaucratic burdock, befuddling binary and digital digamma, that invites us inward toward ever deeper connection with all around us. Yet rather than look within, most in this world prefer to look outward for divine wisdom, thus sustaining a dependence upon a vast hierarchy that stretches between the individual and the Divine. In all our wanderings away from the Divine, humankind has obscured, obfuscated, obliterated, adjured, abnegated and abjured its most compelling features through a persistent belief systems in limitations arising out of the controlling mechanisms of this hierarchy of sub ordinates, superordinates and coordinates. Has my consciousness changed? Why that doubt? Twist and shake, ponderous novelty themes, pathos, poignancy in writing in different styles, accord-coetaneousness, coevality-coincidence, concurrence-conformity, conjunction-contemporaneousness, harmony, order, peace, simultaneousness- synchronicity exciting optimism of youth, power obmbudsman,pulsations and throbbing of pop, hard luck stories of people, modernity and mischief, grimalkins of rurality, wobbling boats of lovers, saxophone in mouths all point to newer value orientations. The Divine dances and the “Grand design “is perpetrated outside of the confines of any hierarchical structures, be it feudal, organizational. It is complete within itself, and has a singular purpose of demonstrating the collective potential of all life within the universe. It is the archetype of perfection. It is the standard bearer of each soul's innate design and ultimate destiny. The essence of the Divine is far beyond mental conception, yet humanity's tendency is to resort to the limiting language of the hierarchical paradigm acceding- accession, accommodation- accord, accordance, acknowledging, adjustment-affiliation, affinity-alliance, amity-approving, arbitration-arrangement, assenting, authorizing, bargaining-compatibility, compliance--compromise, concert-concession, conformity-congruity, consistency-correspondence, endorsing, granting-harmony in conformity to pen chance ,predilections of the protagonist to define and understand it. Searching for Connection and Wholeness calls for not lumping elements like me made of clay feet and weak knees, owing allegiance to the fringe, without courante and courage of conviction and will for vindication, visual restlessness and exaggerated narratives of self to its own, exciting potentialities available, terrestrial passion, and skew of scams all bear an infallible demonstration for new order of consciousness. Theorists have posited a virtual smorgasbord of objects to account for dark matter, although many of them have fallen prey to observational constraints. As leading possible candidates for baryonic dark matter, there are black holes (large and small), brown dwarfs (stars too cold and faint to radiate), sun-size MACHOs, cold gas, dark galaxies and dark clusters, to name only a few. The range of particles that could constitute nonbaryonic dark matter is limited only slightly by theorists' imaginations. The particles include photinos, neutrinos, gravitinos, axions and magnetic monopoles, among many others. Of these, researchers have detected only neutrinos--and whether neutrinos have any mass remains unknown. Experiments are under way to detect other exotic particles. If they exist, and if one has a mass in the correct range, then that particle might pervade the universe and constitute dark matter. But these are very large "ifs." To a great extent, the details of the evolution of galaxies and clusters depend on properties of dark matter. Without knowing those properties, it is difficult to explain how galaxies evolved into the structures observed today. As knowledge of the early universe deepens, I remain optimistic that we will soon know much more about both galaxy formation and dark matter. What we fail to see with our eyes, or detectors, we can occasionally see with our minds, aided by computer graphics. Computers now play a key role in the search for dark matter. Historically, astronomers have focused on observations; now the field has evolved into an experimental science. Today's astronomical experimenters sit neither at lab benches nor at telescopes but at computer terminals. They scrutinize cosmic simulations in which tens of thousands of points, representing stars, gas and dark matter; interact gravitationally over a galaxy's lifetime. A cosmologist can tweak a simulation by adjusting the parameters of dark matter and then watch what happens as virtual galaxies evolve in isolation or in a more realistic, crowded universe.

Old consciousness will not do. Newer one has to be imbibed and cultivated. When people are unaware of their inner wisdom and wholeness, they tend to search for order and security outside themselves. That is my personal experience. There is always search for sanctity of emotion, aesthetic expression, pursuit of harmony, unraveling realms, syncopated ending and stillness of mind and shiver of heart at the mirror broken somewhere are the strains of newer wisdom if it can be called one. Same questions shall be asked and same shall be answer in the similar vein. Uncertain of their place within the hierarchical order of the world, they define themselves based upon their insecurities. Individuals thus become only pieces of their wholeness and like shards of glass from a beautiful vase; they bear little resemblance to their aggregate beauty. Within the hierarchy, many in high positions of power have taken advantage of our collective insecurities in attempting to guide the development of all humankind. They have obscured the direct connection between the individual and the Divine through a variety of means designed to intercede between our inner essence and our divine source. Each individual can come to know them to be free of all forms of institutional and organisational and societal hierarchical control. This is not to imply that we should not trust others or join together in bonds of friendship and community. It is simply a reminder that relative truth is constantly shifting in the hands of those who desire to control. And even when the motive for their controlling behaviors may be of good will, it is still a form of control. When the revealers of "truth" within the hierarchy withhold and suppress information, they are usually positioning themselves to maintain or acquire more power rather than to disseminate empowerment to all. The desire for connection and wholeness is a fuel that drives us to seek out and explore the hierarchical paradigm. This inner longing provides us with the motivation to seek help and guidance from a specific group within the hierarchy, and in so doing, cultivate a sense of belonging and connection. Furthermore, the hierarchical paradigm is a stage whereby we develop a sense of connection to some grand, encompassing vision. This is why the hierarchy nurtures prophets who point toward a greater vision. Spiritual leaders are able to peer deeply beneath the surface reality of life and experience how intricately connected every life form is, and how the composite of all life is intelligent far beyond measure. These visionary leaders can thus interpret reality through their personal abilities to perceive and express life's dimensional depth and limitless intelligence. Yet no one is able to articulate life's full dimensional depth and breadth with the tools of language. They can only, at best, describe their interpretation or their impressions. In actuality, all of us are able to peer beneath the surface reality of life and perceive a unique vision of the universe. We require only time and intention to develop our own interpretations. And this is precisely what many great spiritual leaders have taught. Life's deeper essence is not an absolute to be experienced by the chosen few, but an evolving, dynamic intelligence that wears as many faces as there are life forms. No group or people have the exclusive portal into the universe by which the Divine expresses itself in all its majesty. The portal is open and available to all, because the Divine is within all things. Those recognized as great prophets each produced a vision of the universe beyond what was currently defined by the hierarchy. Because their interpretations were articulated with authority and depth of insight, they became a target of debate among various groups in the hierarchy. This debate then created a polarity of belief. A sympathetic constituency emerged to defend and embellish their leader's interpretation, while established groups held it in contempt of previously held beliefs. Invariably, the leader's vision became confined and shaped into dogma by followers who desired to create a new religion or sect. Thus, this infusion of fresh insight quietly receded into the hands of the hierarchy, where its deeper meaning was obscured by the very fact that it was incorporated into a rigid structure that both protected and promoted it. A new paradigm is emerging that promotes a clear connection of individual consciousness to the compelling features of the Divine without the intervention of a hierarchy. This is when the fables and myths of history step into the light and become known as they were originally intended. This is the time when language will be transformed into a new form of communication that breaks down all barriers of control. Personal transformation, through the awakening of inner wisdom and understanding, is the pathway into wholeness. The transformation paradigm is initiated simply by the recognition that rather than the dependency-inducing ways of the hierarchical paradigm, there are accelerated, independent pathways that bypass the hierarchy and lead to self-mastery. These new pathways lead to the divine wisdom and understanding that is present within all of us. This wisdom can be accessed through the practice of three principles of transformation: seeing the Divine in

all, nurturance of life, and gratitude. The application of these life principles disengages individuals from the controlling elements of the hierarchy, thus initiating the transformation experience. As there are relative truths, there are relative freedoms. As individuals evolve through the hierarchical process, an ever-increasing sense of freedom is gained, yet external forces continue to exert control through limiting language, confining belief systems, and more. These controlling influences lead to anatomical architecture, body, build, cagey carcass, construction, enclosure, fabric, flounce, form, framework, fringe, groundwork, hem, mount, mounting outline, physique, scaffold, scaffolding scheme, setting shell, stage, structure, support, system, trim, trimming, truss- valance continued reliance upon the hierarchy as it unceasingly attempts to impose a sense of inequality between us. The underlying equation of the hierarchical paradigm is: individual + hierarchy = God connection. In the case of personal transformation it is: individual + inner wisdom and understanding divine equality with all. The time has come to integrate the dominant paradigm of the hierarchy with the liberating transformation paradigm. That shall redeem all promises made to oneself. There shall be a Matrix force or balancing force that shall remain UN satisfies until conservativeness is reached. This integration occurs naturally once we have fully explored the two paradigms and develop a synthesis paradigm whereby transformation is attained by searching for connection and wholeness through our own inner wisdom and understanding. It is this combination of release from dependency on the hierarchy and transformation from within that initiates the synthesis paradigm. Once we take responsibility for our own transformation and integration, it does not mean that the hierarchy is to be shunned or avoided. The hierarchy is quite benign as a manipulative force and merely represents one important stage in the journey toward wholeness. What is being set in motion now is the initial preparation for these paradigm shifts. More specifically, these paradigms will be simultaneously played out over the coming years. As always, it will be the choice of each one of us as to which paradigm we embrace in our journey. All of the highest imaginings of humanity are yet unaware of our deepest foundation. We have sought the upper reaches of the building, yet remain unaware of the foundation's design. It is here, at the very core of existence that the Divine is bursting forth with its creative energy while simultaneously reintegrating with its invitation to wholeness. It is here that equality is realized, not in the lofty places of relative truth lodged in the hierarchy, but rather in the deepest part of the foundational plan of life's original source and ultimate destiny. The origin and destiny of existence is the tone of equality in life. Listen for this tone – this frequency of vibration – and follow it back into the very foundation from whence all things arise and return. Consider these words as symbols only. Feel the truth that stands behind these words, and tap into this empowering energy force that reaches out for you. Know it as a tone or vibration – a resonance that waits for you around every corner in which your life will turn. It is a beacon of the Divine gathering itself into the form of language in order to usher you to a place from which you can experience the formless tone of equality – the bypass of limitation. It is the primal language of our divine source that bestows to you the freedom to generate your own deepest beauty in the expression of your highest truth.(Reference: Wikipedia)

## Gopalasamy Model For Procrastinated Conducive, Felicitous Counter Sallies In Two Specified Systems

Abstract

Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of decay rate of perturbations. Sole intention, main objective and primary concern of

the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are counted in the same vein as that of predator prey. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 5. Preliminary remarks

Almost all the equations of population ecology have been derived starting from the following simple format (Note this can be used for the development of the creed and type of systemic types in case of cosmology and astrophysics and fiancé and econometrics. To give an example of finance, we can think of a account of the company giving rise to various accounts of its employees, and that is the contribution of the single account for the development of the total system types of the given configuration). It may be account, particle, atom, molecule, proposition, problem or any such thing. We shall delineate and disseminate information and particulars as and when the model shall be applied to the problem in question.

$$\frac{dN(t)}{dt} = \text{an individual's contribution}$$

*to population growth* $N(t)$

(1.1)

where $N(t)$ denotes the size of a population of a single species (or biomass) with overlapping generations at time $t$, and $dN(t)/dt$ is the rate of change of the population size at time $t$. Subsequently one makes an assumption regarding the factor in the brackets in (1.1); in particular if one assumes that an individual entity's contribution to the population growth is denoted by a function say, f(t, N), defined suitably for all t>0, N then one obtains from (1.1) the so called Kolmogorov formulation in the form

$$\frac{dN(t)}{dt} = f(t, N(t))\, N(t).$$

(1.2)

Various choices of $f$ together with some ecologically plausible assumptions such as the temporal constancy of the environment and density dependent effects in $f$ lead to various equations. For instance if $f(t, N) \equiv r$ (a positive constant) one obtains the Malthusian formulation in the form

$$\frac{dN(t)}{dt} = rN(t)$$

And if one assumes $f(t, N) \equiv r - (r/K) N$ for some positive constants r and K one gets the familiar logistic equation of population biology

$$\frac{dN(t)}{dt} = rN(t) \left\{ 1 - \frac{N(t)}{K} \right\}$$

In order to device models with time delayed effects, such effects are introduced in the average growth rate $F(t, N)$ of (1.2); for instance if $f(t, N) \equiv r - (r/K) N(t - \tau)$ for some constant $\tau > 0$ we obtain Hutchinson's[19]delay loistic equation

$$\left\{ \frac{dN(t)}{dt} = rN(t) \left\{ 1 - \frac{N(t-\tau)}{K} \right\} \right. .$$

(1.3)

If for instance

$$F(t, N) = r - (r/K) N(t) - \int_{-\infty}^{t} k(t - s) F(N(s))\, ds$$

Where $k$ is a piecewise continuous nonnegative valued function such that $\int_0^\infty k(s)\,ds = 1$ with some appropriate function $F$, we will get Volterra's model of a population which pollutes its environment and pollution has an accumulative and toxic effect; such an equation is of the form

$$\frac{dN(t)}{dt} = N(t) = \left\{ r - (r/K) N(t) - \int_{-\infty}^{t} k(t-s) F(N(s))\,ds \right\} \quad (1.4)$$

Which has been studied in different formats by several authors (Cushing [12], Cohen *et al.* [10], and Gopalasamy [16])

One of the significant advantages of deriving various models from the prototype (1.2) is that if the initial; conditions (initial population sizes) are nonnegative then the nonnegativity of the population size $N(t)$ for $t > 0$ follows from the fact that solutions of (1.2) satisfy

$$N(t) = N(0) \exp\left[\int_0^t f(s, N(s))ds\right]$$

Which implies that $N(t) \geq 0$ if $N(0) \geq 0$.

While the time delayed logistic equation (1.3) explains some of the observed oscillations of certain populations, its derivation has not been entirely satisfactory. For instance it has been pointed out by Ricklefs ([27], page 488) that "time lags might be expected to occur primarily in stabilizing processes involving production rather than death because death is an immediate response to environmental change". Recently Cushing [13] has examined the effects of response and maturation delays in a number of single and multi system models; as it has been noted by Cushing [13], the models involving delayed reproduction or not necessarily derivable from the Kolmogorov type systems of the form (1.2) by suitably defining the average growth rate $(1/N)(dN/dt) = f(t, N(t))$. While Cushing [13] has been mostly concerned with the possibility of delay induced oscillations we will consider a class of delays in reproduction and interspecific interaction which maintain the stability of the systems.

Our starting point to drive models with delayed reproduction is the following balance equation which assumes that there is no immigration or emigration. Account, particle, atom, molecule, proposition, problem et al

$$\frac{dN(t)}{dt} = \text{birth rate} - \text{death rate}. \quad (1.5)$$

Now for instance if we consider a population of adult flies then the birth rate at time $t$ depends on the adult population at time $t^{-}\ \tau$ where $\tau$ is the time required for the larvae to become adult flies. If the birth and death rates are governed by density dependent factors then we have from (1.5) the following time delayed model

$$\frac{dN(t)}{dt} = b\big(N(t^{-}\tau)\big) - m(N(t))$$

(1.6)

Where the functions $b$ and $m$ denote the density dependent birth and death rates of Account, particle, atom, molecule, proposition, problem et al respectively. If the time delay is continuously distributed then one can consider in the place of (1.6) an equation of the form

$$\frac{dN(t)}{dt} = b\int_{-\infty}^{t} k\{(t-s)N(s)ds - m(N(t))$$

Or equivalently

$$\left\{\frac{dN(t)}{dt} = b\int_{0}^{\infty} k(s)\right\}N(t-s)ds - m(N(t))$$

(1.7)

Where the delay kernel $k$ denotes the weight of the delayed of the delayed effects.

It is not always obvious that suitable nonnegative initial conditions for (1.6) and (1.7) will imply the nonnegativity of the solutions of (1.6) and (1.7). The purpose of this article is to show that in two species systems of competitive and prey- predator communities, delayed reproduction and delayed interactions donot destabilise the systems if the self- regulating negative feedback effects are sufficiently strong compared to interspecific interactive effects and if the self regulating effects are realised with no time delays. When we say that time delays do not destabilise, we will mean that delays do not render an asymptotically (locally) stable steady state unstable in such a manner that the loss of stability either leads

To a delay – induced bifurcation to persistent and undamped oscillations (as in the case of Hopf- bifurcation) or the relevant linear variational system has unbounded solutions; that is a mathematical instability is not induced by delays in the sense, that the roots of the characteristic equation are pushed to the imaginary axis or to the right of it on the complex plane. However delays can make an otherwise stable system less stable which means that the rate of convergence to zero of solutions of the variational system can decrease with an increase of delay and the variational system can even oscillate (perhaps violently) before converging to the trivial state. Such a reduction in stability can

happen even when all the roots of the characteristic equation have negative real parts. For our purposes such a reduced stability is still stability. In order to see the effects of delays on such less stable systems we will propose below a method of estimating the rate of decay of local perturbations based on the variational system and indirectly relying on the characteristic equation. For a special class of two species models with a single delay we have obtained (Gopalsamy [16]) estimates of decay rates (or return times) usingthe characteristic equation; such a procedure is almost impossible in the case of many species systems with several delays. The method we propose in Section 6 below for the estimate of decay rates can be adapted to *n*- species systems with $n^2$ discrete delays.

Let $x_1(t)$ and $x_2(t)$ denote the population sizes (or biomasses) of two systems competing for a common pool of resources.. Let $b_i$ and $m_i$ $(i = 1, 2)$ denote the respective density dependent birth (potential) and death rates so that in the absence of any time delays the population sizes are governed by (of Account, particle, atom, molecule, proposition, problem et al). In addition as usual, we assume that the system studied follow Haimovici's equations, for which the dove tailing explanation has been appended

$$\left.\begin{aligned}\frac{dx_1(t)}{dt} &= b_1\left(x_1(t)\right) - m_1(x_1(t)), x_2(t))\\ \frac{dx_2(t)}{dt} &= b_2\left(x_2(t)\right) - m_2(x_1(t)), x_2(t))\end{aligned}\right\} \quad (2.1)$$

In order to make the system (2.1) represent a competition system we will make the following assumptions on the birth and death rates; $b_1$, $m_1$ $(i = 1, 2)$ are continuous with continuous partial derivatives for all $x_1 \geq 0$ $(I = 1, 2)$; also we assume (see for instance Rescigno and Richardson [26]) $\frac{\partial b_i}{\partial x_1} > 0, \frac{\partial m_1}{\partial x_j} > 0$ for $x_i > 0$ $(i, j = 1, 2)$, (2.2)

(vii) $b_1(0) = 0$, $m_1(0, x_2) \equiv 0, b_2(0) = 0$, $m_2(x_1, 0) \equiv 0$, (2.3)For some $x_{1,}^* > 0, x_{2,}^* > 0$ we have

$$\left.\begin{aligned}b_1(x_1^*, -m_1(x_1^*, 0) &= 0,\\ b_2(x_2^*, -m_1(0, x_2^*) &= 0,\end{aligned}\right\} \quad (2.4)$$

(viii) There exist positive constants $\delta_1$, $\delta_2$ such that

$b_1(\delta_1) - m_1(\delta_1, x_2) < 0$,

$b_2(\delta_2) - m_2(x_1, \delta_2) < 0$, for $x_1 \geq 0$, $x_2 \geq 0$,

(2.5)

(ix) There exists a point $(\alpha, \beta), \alpha > 0, \beta > 0$ such that

$$b_1(\alpha) - m_1(\alpha, \beta) = 0 = b_2(\beta) - m_2(\alpha, \beta).$$

(2.6)

The conditions on $b_i$ in (2.2) mean that the potential (not average) birth rates are density dependent and any negative density dependent feedbacks to the opened or birth rate of Account, particle, atom, molecule, proposition, problem et al rates are included in the destruction rates; the conditions on $m_i$ indicate intraspecific and interspecific confederation needs and necessities. (2.3) implies that (0, 0) is a (trivial) steady state of the system (2.1) as it is in the case of all ecological models; (2.4) means that in the absence of either of the two species the other one has a nontrivial steady state; (2.5) implies that each of the two populations cannot grow unboundedly since when $x_1(t) > \delta_i, \frac{dx_1}{dt} < 0$. (2.6) means that there exists a nontrivial steady state $(\alpha, \beta)$ in the interior of the positive quadrant of the $(x_1, x_2)$- plane. These assumptions are standard and are minimal in the opinion of the author.

An immediate example of (2.1) –(2.6) is the familiar Volterra –Lotka competition system described by

$$\frac{dx_1(t)}{dt} = r_1 x_1(t) - a_{11} x_1^2(t) - a_{12} x_1(t) x_2(t),$$

$$\frac{dx_2(t)}{dt} = r_2 x_2(t) - a_{21} x_1(t) x_2(t) - a_{22} x_2^2(t),$$

(2.7)

Where $r$, $a_{ij}$, $(i, j = 1,2)$ are positive constants satisfying the condition

$$\frac{a_{11}}{a_{21}} > \frac{r_1}{r_2} > \frac{a_{12}}{a_{22}}.$$

One of the first questions to be resolved for (2.1) is the following: if $x_i(0) > 0$, $x_2(0) > 0$ does it follow that the corresponding solutions of (2.1) satisfy the condition $x_1(t) \geq 0$, $x_2(t) \geq 0$ for those $t$ for which a solution of (2.1) exists? In the $(x_1, x_2)$ phase plane of (2.1) we have (0, 0),$(x_1^*, 0)$,$(0, x_2^*)$, $(\alpha, \beta)$ as the steady states of (2.1) and also it will follow from (2.3) that the co ordinate axes

$x_1 = 0$ and $x_2 = 0$ are invariant sets f (2.1) ; such invariance of the coordinate axes will imply that if $x_1(0) > 0$, $x_2(0) > 0$ then the corresponding population trajectory $(x_1(t), x_2(t))$ of (2.1) cannot reach the outside of the nonnegative quadrant $x_1 \geq 0$, $x_2 \geq 0$.

Our plan is to derive first a set of sufficient conditions for the coexistence of the two competing species and then in the next section show that under the same conditions the two species can coexist if there are time delays in the reproduction and interspecific interactions. We recall that the two competing species described by (2.1) are said to coexist if and only if the nontrivial steady state $(\alpha, \beta)$ is (at least locally) asymptotically stable. To examine such a stability of $(\alpha, \beta)$ we derive the corresponding linear variational system for the perturbations $X_1, X_2$ where

$$x_1(t) \equiv \alpha + X_1(t), \qquad x_2(t) \equiv \beta + X_2(t)$$

Such a variational system is found to be

$$\left.\begin{aligned} \frac{dX_1(t)}{dt} &= \beta_{11} X_1(t) + \beta_{12} X_2(t), \\ \frac{dX_2(t)}{dt} &= \beta_{21} X_1(t) + \beta_{22} X_2(t), \end{aligned}\right\} \qquad (2.8)$$

Where

$$\beta_{1J} = \left(\frac{\partial b_1}{\partial x_J} - \frac{\partial m_1}{\partial x_j}\right) \text{ evaluated at } (\alpha, \beta) \qquad (i, j = 1,2).$$

(2.9)

The steady state $(\alpha, \beta)$ of (2.1) is locally asymptotically stable if and only if the trivial solution (0, 0) of (2.8) is asymptotically stable. We can derive the following result from elementary considerations.

In the system (2.1)-(2.6) assume the following:

$$\left.\begin{aligned} \frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1}, \\ \frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1}, \end{aligned}\right\} \qquad \text{at } (\alpha, \beta)$$

(2.10)

Then the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (2.1) – (2.6).

The proof follows easily if we note that the characteristic equation associated with the linear system (2.8) is given by

Det $\begin{pmatrix} \lambda-\beta_{11} & -\beta_{12} \\ -\beta_{21} & \lambda-\beta_{22} \end{pmatrix} = 0$

Or equivalently,

$$\lambda^2 - \lambda\,(\beta_{11} + \beta_{22}) + \beta_{11}\beta_{22} - \beta_{12}\beta_{21} = 0. \qquad (2.11)$$

It follows from (2.2) and (2.10) that,

$$\beta_{11} < 0, \beta_{11}\beta_{22} - \beta_{12}\beta_{21} > 0$$

And hence the roots of (2.11) have negative real parts implying the asymptotic stability of the trivial solution of (2.8)

If apply the condition (2.10) to the Volterra-Lotka model (2.7) we note that when a nontrivial steady state say $(\alpha, \beta)$ exists for (2.7), (2.10) leads to

$$a_{11}\alpha > a_{21}\,\beta \; and \; a_{22}\,\beta > a_{12}\,\alpha,$$

Which together imply $a_{11}\,a_{22} > a_{12}a_{21}$. If $(\alpha, \beta)$ exists it is known that $a_{11}\,a_{22} > a_{12}a_{21} > 0$ is a sufficient condition for the (local) asymptotic stability of $(\alpha, \beta)$.

The condition (2.10) which guarantees the coexistence of competing populations in (2.1)-(2.6) simply means that the intraspecific self – regulating feedback effects $\partial m_i / \partial x_1\ (i = 1, 2)$ bare higher than its own positive feedback $\partial b_1 / \partial x_1\ (i = 1, 2)$ together with its negative (competitive) effects on its competitor $\partial m_I / \partial x_J\ (i,\ j = 1, 2; i \neq j)$.

## 6. Coexistence under delayed production and interactions, transactions

We will now establish that two competing species with a sufficiently strong intraspecific regulation can coexist in the presence of delayed reproduction and delayed interspecific interactions. It is well known that if there are delays in the self – regulating intraspecific effects, stationary coexistence may not always be possible (Gopalsamy [16], Gopalsamy and aggarwala [15], Shibata and Saito [28]). Delayed reproduction is more general and widespread in most biological populations. To build a competition model with such delays we proceed as follows.

Let $\tau_{ij}\ (i, j = 1, 2)$ be a set of nonnegative constants with $\tau = \max\ \tau_{ij}\ , i, j = 1, 2,$ and suppose that the two ompeting species display delayed reproduction and delayed interspecific interactions with no delays in the self – regulating negative feedback effects. Such a competition system in a constant environment can be modeled by the autonomous delay differential system

$$\frac{dx_1(t)}{dt} = b_1\big(x_1(t - \tau_{11})\big) - m_1\big(x_1(t), x_2(t - \tau_{12})\big),$$

$$\frac{dx_2(t)}{dt} = b_2\big(x_2(t - \tau_{22})\big) - m\big(x_1(t - \tau_{21}), x_2(t)\big), \qquad (3.1)$$

Where both the creation and destruction rates and the morality rates $b_1$ and $m_1\ (i = 1, 2)$ satisfy the same conditions (2.2)-(2.6) as in the previous section. Along with (3.1) we will

Suppose that the initial population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) sizes are provided in the form

$$x_i(s) = \phi_i(s) \geq 0, \qquad s \in [-\tau, 0], \qquad i = 1,2. \tag{3.2}$$

(Note that the specification of $\phi_i$ need not be on the entire interval$[-\tau, 0]$; but such a hypothesis is convenient.) Since (3.1) is not of the standard Kolmogorov-type, we have to establish that the solutions of (3.1) will remain nonnegative for nonnegative initial conditions.

To prove the nonnegativity of solutions of (3.1)-(3.2), let $\left\{\psi_1^{(n)}(s)\right\}, s \in [-\tau, 0], i = 1,2; (n = 1,2, \ldots)$ be a sequence of strictly positive continuous functions such that

$$\lim_{n\to\infty} \psi_1^{(n)}(s), = \psi_1(s), \qquad s \in [-\tau, 0], \qquad i = 1,2 \tag{3.3}$$

Let $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ be the solution of (3.1) corresponding to the initial condition

$$\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \left\{\psi_1^{(n)}(s), \psi_2^{(n)}(s)\right\}, \quad s \in [-\tau, 0],. \tag{3.4}$$

Consider the solution $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ for $t \in [0, \tau]$ (assuming it exists on the whole of $[0, \tau]$). Suppose $x_1^{(n)}$ does not remain nonnegative for all $[0, \tau^*]$; then there exists a $t^* = [0, \tau^*](\tau^* =$ the positive minimum of $\tau_{ij}$ $(i, j = 1, 2))$ for which

$$x_1^{(n)}(t^*) = 0 \ and \ x_1^{(n)}(t)(t^* - \tau_{11})) > 0.$$

It will follow from (3.1) and the positivity of the initial condition that

$$\frac{dx_1^{(n)}(t^*)}{dt} = b_2\left(x_1^{(n)}(t^* - \tau_{11})\right) > 0 \tag{3.5}$$

Where we have used the properties of $b_i$ and $m_i$ from (2.2)-(2.6); it is found that (3.5) contradicts the definition of $t^*$. Thus $x_1^{(n)}(t) > 0$ on $[0, \tau^*]$; similarly $x_2^{(n)}(t) > 0$ $on$ $[0, \tau^*]$. We can repeat this procedure for intervals of the form $[\tau^*, 2\tau^*], [2\tau^*, 3\tau^*]$ $etc$. Thus it will follow that so long as $(x_1^{(n)}(t), x_2^{(n)}(t))$ exists, we have

$$x_1^{(n)}(t) \geq 0 \, x_2^{(n)}(t) \geq 0.$$

If we now consider the limits as $n \to \infty$, we get

$$\lim_{n\to\infty}\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \{x_1(t), x_2(t)\} \tag{3.2}$$

With $x_1(t) \geq 0, , x_2(t) \geq 0$ where $(x_1(t), x_2(t))$ is the solution of (3.1)-(3.2) and this is a consequence of the continuous dependence of solutions of (3.1) on initial conditions (Hale [17], page 41).

A second question for (3.1) with (2.2)-(2.6) is whether solutions of (3.1) exist for all $t > 0$? suppose now that solutions of (3.1) do not exist for all $t$ ; then there exists a $t_1 > 0$ such that at least for one of either of $x_1$ *or* $x_2$ we have

$$\lim_{t \to t_1 -} x_i\ (t) = \infty. \tag{3.7}$$

To be specific let us suppose that $\lim_{t \to t_1 -} x_i\ (t) = \infty$; let $t_2$ be the first instant of time for which $x_1(t_2) = \delta_1, t_2 < t_1$. It will follow from (3.1) that

$$\frac{dx_1(t_2)}{dt} = b_1\big(x_1\ (t_2 - \tau_{11})\big) - m_1\big(x_1(t_2), x_2\ (t_2 - \tau_{12})\big)$$

$$< b_1\big(x_1\ (t_2)\big) - m_1\big(x_1(t_2), x_2\ (t_2 - \tau_{12})\big)$$

$$< b_1(\delta_1) - m_1\big(\delta_1, x\ (t_2 - \tau_{12})\big) < 0 \tag{3.8}$$

Which contradicts the definition of $t_2$. Thus solutions of (3.1), (2.2) – (2.6) exist for all $t$ .

We are now ready to consider the effects of time delays in reproduction and interspecific interactions on the coexistence of the two competing species. The steady state $(\alpha, \beta)$ of (2.1) is again a steady state of (3.1). To examine the (local) asymptotic stability of $(\alpha, \beta)$ in (3.1) we let

$$x_1(t) = \alpha + Y_1(t), \quad x_2(t) = \beta + Y_2(t)$$

In (3.1) and derive that the linear variational system governing the perturbations $Y_1, Y_2$ is given by

$$\frac{dY_1\ (t)}{dt} = -m_{11}Y_1(t) + b_{11}Y_1\ (t - \tau_{11}) - m_{12}Y_2\ (t - \tau_{12}),$$

$$\frac{dY_2\ (t)}{dt} = -m_{22}Y_2(t) + b_{21}Y_1\ (t - \tau_{21}) - m_{22}Y_2\ (t - \tau_{22}),$$

Where

$$m_{ij} = \frac{\partial m_1}{\partial x_j}, \quad b_{ii} = \frac{\partial b_i}{\partial x_i}\ at\ (\alpha, \beta). \tag{3.10}$$

It is well known (Bellman and Cooke [3], page 336) that the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) if the trivial solution of the variational system (3.9) is asymptotically stable. We formulate our first result in the following form whose proof is given in the appendix.

*THEOREM 3.1. Assume that in the two systemic types competition model system (3.1) with delayed density dependent birth and death rates $b_i$ and $m_i$ $(i = 1, 2)$ satisfy the conditions (2.2)-(2.6). Suppose that the conditions (2.10) hold. Then the nontrivial steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) whatever the delays be $\tau_{ij}$ $(i = 1, 2)$.*

Let $y_1(t) and\ y_2(t)$ respectively denote the population sizes (or biomasses) of a prey species and a predator species which feeds exclusively on the single prey species in a temporally constant environment. In the format of (1.5) the following system will describe the rates of change of the two population sizes:

$$\frac{dy_1(t)}{dt} = f_1(y_1(t)) - g_1(y_1(t), y_2(t)),$$
$$\frac{dy_2(t)}{dt} = f_2(y_1(t)) - y_2(t) - g_2(y_2(t)), \quad (4.1)$$

Where $f_i\ and\ g_i\ i = 1,2)$ respectively denote the birth and death rates of the two species; in order to make (4.1) denote the dynamics of a prey – predator species we will assume that $f_i\ and\ g_i$ satisfy the following conditions (see for instance Rescigno and Richardson [26]); $f_i\ and\ g_i$ are continuous with continuous partial derivatives for all $y_i \geq 0\ (i = 1, 2); also\ we\ require$

$$\frac{\partial f_1}{\partial y_1} > 0; \quad \frac{\partial g_1}{\partial y_1} \geq 0; \quad \frac{\partial g_1}{\partial y_2} \geq 0,$$
$$\frac{\partial f_2}{\partial y_1} \geq 0; \quad \frac{\partial g_2}{\partial y_2} \geq 0; \quad \frac{\partial g_2}{\partial y_2} \geq 0, \quad \text{Where } y_1 \geq 0, y_1 \geq 0, \quad (4.2)$$

(ix)

$$f_1(0) = 0, \quad g_1(0, y_2 \equiv 0,$$
$$f_2(0, y_2) = 0 = f_2(y_1, 0); \quad g_2(0) = 0, \quad (4.3)$$

(x) For some constant $y_1^* > 0$ we have

$$f_1(y_1^*) - g_1(y_1^* 0) = 0, \quad (4.4)$$

(xi) There exists a point $(\alpha^*\beta^*), \alpha^* > 0, \beta^* > 0$ such that

$$f_1(\alpha^*) - g_1(\alpha^*, \beta^*) = 0,$$
$$f_2(\alpha^*, \beta^*) - g_2(\beta^*) = 0, \quad (4.5)$$

(xii) There exist positive constants $n_1, n_2$ such that

$$f_1(\eta_1) - g_1(\eta_1, y_2) < 0, y_2 \geq 0,$$
$$f_2(\eta_1, \eta_2) - g_2(\eta_2) < 0 \quad (4.6)$$

The conditions (4.2) mean that the system (4.1) is of the oppugnacy systems type; (4.3) will imply that (0,0) is a trivial steady state of (4.1); (4.4) means that in the absence of the predator, the prey species has a nontrivial (positive) steady state;

(4.5) implies that (4.1) has a nontrivial steady state in the interior of the positive quadrant of the state species of (4.1); (4.6) can be shown (see below) to imply that no credit type and debit type system types can grow unboundedly.

An example of (4.1)-(4.6) is the familiar Volterra-Lotka prey-predator system

$$\frac{dy_l(t)}{dt} = r_l y_l(t) - [a_{ll} y_l^2(t) + a_{12} y_l(t) y_2(t)],$$
$$\frac{dy_2(t)}{dt} = a_{12} y_1(t) y_2(t) - [r_2 y_2(t) + a_{22} y_2^2(t)],$$

Where $r_{l,}\ a_{ij}\ (i,j=1,2)$are positive costants such that

$$a_{21}r_1 - r_2a_{11} > 0.$$

As in the previous case of competition we have to verify that when $y_1(0) > 0, y_2(0) > 0$ the corresponding solution of (4.1) will be nonnegative for $t \geq 0$ for which the solution exist. Since the co-ordinate axes $y_1 = 0, y_2 = 0$ are invariant sets for (401) by virtue of (4.3), it will follow that the solution of (4.1) starting from the interior of the positive quadrant cannot reach the outside of that quadrant. This is a consequence of the invariance of the co-ordinate axes for (4.1)

To prove that solutions of (4.1)-(4.6) exist for all $t \geq 0$ let us suppose that $y_l 0\ and\ y_2(0) satisfy\ 0 < y_l\ (0)$ satisfy $0< y_l(0) < \eta_l, 0 < y_2(0) < \eta_2$ ( on the other hand if either one of $y_l(0) \geq \eta_l and\ y_2(0) \geq \eta_l$ or both hold, it will follow that the corresponding $y_l(t)$ will decrese by virtue of (4.6) at least for small $t > 0,$ ). Suppose $y_l(t)$ does not exist for all $t \geq 0$; then for some $t_l < \infty\ lim_{t\to t_l-} y_l(t) = \infty$and let $t_2 > 0$ be the first instant for which $y_1(t_2) = \eta_1$; then we have from (4.1) and (4.6) that

$$\frac{dy_1(t_2)}{dt} = f_1(\eta_1) - g_1(\eta_{1,}\ y_2) < 0\ (assumimg\ y_1(0) < \eta_1)$$

Which contradicts the definition of $t_2$ showing that $y_1(t)$ exists for all $t \geq 0$; if $y_2(t)$ is not defined for all $t \geq 0$, then for some $t_3 \geq \infty$ we will have $lim_{t-t_3-} y_2(t) = \infty$;now let $t_4$ be the first instant for which $y_2(t_4) = \eta_2$; then from the properties of $f_2$ we have

$$\frac{dy_2(t_4)}{dt} = f_2(y_1(t_4), \eta_2) - g_2(\eta_2) \leq f_2(\eta_1, \eta_2) - g_2(\eta_2)$$

$$< 0 (by\ (4.6) assuming\ 0 < y_1(0) < \eta_1, 0 < y_2(0) < \eta_2)$$

Which again contradicts the definition of $t_4$. Thus $y_2(t)$ is defined for all $t \geq 0$.

We will now derive a set of sufficient condition for the (local) asymptotic stability of the steady state $(\alpha^*, \beta^*)$ of (4.1). in the next section we will show that under the same conditions and in the presence of various delays , the steady state $(\alpha^*, \beta^*)$remains (locally) asymptotically stable.

To examine the asymptotic stability of the nontrivial steady state $(\alpha^*, \beta^*)$ in (4.1) we let

$$y_1(t) = \alpha^* + U_1(t), \qquad y_2(t) = \beta^* + U_2(t)$$

In (4.1) and derive the liner variational system in the perturbations $U_1, U_2$ as follows:

$$\left.\begin{aligned} \frac{dU_1(t)}{dt} &= \frac{\partial f_1}{\partial y_1}U_1(t) - \frac{\partial g_1}{\partial y_1}U_1(t) - \frac{\partial g_1}{\partial y_2}U_2(t), \\ \frac{dU_2(t)}{dt} &= \frac{\partial f_2}{\partial y_1}U_1(t) - \frac{\partial f_2}{\partial y_2}U_2(t) - \frac{\partial g_2}{\partial y_2}U_2(t), \end{aligned}\right\} \quad (4.7)$$

Where all the partial derivatives in (4.7) are evaluated at $(\alpha^*, \beta^*)$. The steady state $(\alpha^*, \beta^*)$ of (4.1) is (locally) asymptotically stable if and only if the trivial solution of (4.7) is asymptotically stable which will be the case if and only if the roots of the characteristic equation

$$\text{Det} \begin{pmatrix} \lambda - (f_{11} - g_{11}) & g_{12} \\ -f_{21} & \lambda - (f_{22} - g_{22}) \end{pmatrix} = 0,$$

Where $f_{ij} = \frac{\partial f_1}{\partial y_l}$, $g_{ij} = \frac{\partial g_1}{\partial y_j}$, at $(\alpha^*, \beta^*)$ (4.8)

$(i,j = 1, 2)$

Have negative real parts. If $\lambda_1, \lambda_2$ are the riits of (4.8) then

$$\lambda_1 + \lambda_2 = (f_{11} - g_{11}) + (f_{22} - g_{22}), \quad (4.9)$$

$$\lambda_1 \lambda_2 = (f_{11} - g_{11})(f_{22} - g_{22}) + f_{21} g_{12}.$$

The proof of the following result is an immediate consequence of (4.7)-(4.9)

THEOREM 4.1 *In the CREDIT-DEBIT system (4.1)-(4.6) assumes that the partial derivatives satisfy the following additional conditions:*

$$\frac{\partial g_1}{\partial y_1} > \frac{\partial f_1}{\partial y_1} + \frac{\partial f_2}{\partial y_1}, \quad \frac{\partial g_2}{\partial y_2} > \frac{\partial f_2}{\partial y_2} + \frac{\partial g_1}{\partial y_2}, \quad \text{Evaluated at } (\alpha^*, \beta^*), \quad (4.10)$$

*Then the steady state $(\alpha^*, \beta^*)$, is (locally) asymptotically stable for the prey- predator system.*

## 7. Delays in reproduction, predation and gestation

A number of authors have proposed and analysed several models of prey- predator systems with time delays. Most of the models have been derived from the so called models of the Kolmogorov-type f (1.2) with time delays incorporated in the average growth rates of the prey and predator species. Volterra [32] proposed the following model for a prey-predator system. The equations are used for the formulation of Credit-debit system of atoms, accounts, economic variable, cosmological variable, astrophysical; variables, and qubit systems for which the model is being applied.

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \delta_1 N_1(t) - \int_{\infty}^{t} F_1(t-s) N_2(s) ds],$$

(5.1)

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 + \delta_2 N_1(t) - \int_{\infty}^{t} F_2(t-s)N_1(s)ds],$$

Where $r_1, r_2, \delta_1, \delta_2$ *are positive constant and* $F_1, F_2$ is nonnegative continuous delay kernel defined and integrable on $[0, \infty]$. (5.1) has no predator self-regulating negative feedback. Brelot [6] considered a modified format of (5.1) in the form

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \lambda_1 N_1(t) - \delta_1 N_2(t) - \int_{-\infty}^{t} F_1(t-s)N_2(s)ds],$$

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 - \lambda_2 N_2(t) - \delta_2 N_1(t) - \int_{-\infty}^{t} F_2(t-s)N_1(s)ds], \quad (5.2)$$

Which has term representing intraspecific competition among both predator and prey species. Under suitable condition on the various parameters in (5.2) one can show that (5.2) has a (locally) asymptotically stable steady state. A number of intergrodifferential equation models of the types (5.1) and (5.2) have been investigated by Cushing [12]

Starting from Hutchinson's (1948)[19] delay logistic equation, May[23] has proposed the following system

$$\frac{dN_1(t)}{dt} = rN_1(t)\left[1 - \frac{N_1(t-\tau)}{K}\right] - \alpha N_1(t)N_2(t),$$

$$\frac{dN_2(t)}{dt} = -bN_2(t) + \beta N_2(t)N_1(t), \quad (5.3)$$

Where r, $\tau$, K, $\alpha, \beta, b$ are positive constants; (5.3) contains a single discrete delay; one can modify (5.3) and incorporate a continuously distributed delay in it so that (5.3) can be written as (May [23])

$$\frac{dN_1(t)}{dt} = rN_1(t)[1 - \int_{-\infty}^{t} F(t-s)N_1(s)ds] - \alpha N_1(t)N_2(t),$$

$$\frac{dN_2(t)}{dt} = -bN_2(t) + \beta N_1(t)N_2(t). \quad (5.4)$$

The first model of a prey-predator type which departs from the Kolmogorov-type formulation is due to Wangersky and Cunningham [33] who considered the model system

$$\frac{dN_1(t)}{dt} = a_1N_1(t) - b_1N_1^2(t) - c_1N_1(t) - N_2(t),$$

$$\frac{dN_2(t)}{dt} = a_2N_2(t) + c_2N_1(t-\tau)N_2(t-\tau), \quad (5.5)$$

Where $a_1, a_2, b_1, c_1, c_2, \tau$ are positive constants; (5.5) means that a duration of $\tau$ time units elapses when an individual prey is killed and the moment when the corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) Recently Cushing [13] has considered a model of the form

$$\frac{dN_1(t)}{dt} = rN_1(t)\left[1 - \frac{N_1(t)}{k} - \alpha N_2(t)\right],$$

$$\frac{dN_2(t)}{dt} = -d_2^*(t) + bN_1(t)\int_0^\infty \delta(a)N_2(t-a)e^{-d*a}\,da,$$

(5.6)

Where $r, k, \alpha, d^*, b$ are positive constants and $\delta$ is related to the age dependent fecundity of the predator species.

With the models (5.1)-(5.6) in the background let us now consider a time delayed model similar to (4.1) in the following form:

$$\frac{dN_1(t)}{dt} = f_1\big(N_1(t-\tau_{11})\big) - g_1\big(N_1(t), N_2(t-\tau_{12})\big),$$

$$\frac{dN_2(t)}{dt} = f_2\big(N_1(t-\tau_{21})\big), N_2(t-\tau_{12}) - g_1(N_2(t)),$$

(5.7)

Where $\tau_{ij}$ $(i,j=1,2)$ are nonnegative constants with $\tau = \max \tau_{ij}, i,j = 1,2$; the birth and death rate functions $f_1, f_2, g_1, g_2$ satisfy the same conditions as in (4.1)-(4.6). Along with (5.7) we specify the initial conditions as follows:

$$N_1(s) = \phi_1(s) > 0,$$

$$s \in [-\tau, 0].$$

$$N_1(s) = \phi_1(s) > 0,$$

We3 will assume that in addition to (4.2) we have for (5.7) that

$$\frac{\partial f_2}{\partial N_1} > 0, \frac{\partial f_2}{\partial N_2} > 0, f_2(N_2, N_2) > 0 \; whenever \; N_1 > 0, \; N_2 > 0. \; (5.9)$$

It is not difficult to show by the methods previously used for the competition system that whenever $N_1(s) > 0$ and $N_2(s) > 0$ for $s \in [-\tau, 0]$, solution of (5.7) exist for all $t \geq 0$ and remain nonnegative for all $t \geq 0$. also $(\alpha^*, \beta^*)$ of (4.5) is a steady state of (5.7).

We will discuss the form and meaning of the delays in (5.7) in the next section. The effects of the various delays in (5.7) on the (local) asymptotic stability of $(\alpha^*, \beta^*)$ can be derived from the following theorem whose proof is quite similar to that of Theorem 3.1 and hence the details of the proof are omitted.

THEOREM 5.1 *In the time delayed prey-predator system(* corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) *) (5.7) let the conditions (4.2)-(4.6), (5.8) and (5.9) hold for the birth and death rates; furthermore suppose the condition (4.10) hold. Then for all nonnegative delays in (5.7), its nontrivial steady state $(\alpha^*, \beta^*)$ is (locally) asymptotically stable.*

## 6. Estimation of decay rates

We will first derive a convention sufficient condition for the asymptotic stability of the trivial solution of a linear system of the form

$$\left.\begin{aligned} \frac{dY_1(t)}{dt} &= -m_{11} Y_1(t) + b_{11} Y_1(t-\tau_{11}) - m_{12} Y_2(t-\tau_{12}), \\ \frac{dY_2(t)}{dt} &= -m_{22} Y_1(t) + m_{21} Y_1(t-\tau_{21}) - b_{22} Y_2(t-\tau_{22}), \end{aligned}\right\} \; t > 0, \quad (6.1)$$

Where $m_{1J}, b_{i1}$ $(i, j = 1,2)$ are positive constants while $\tau_{1J}(i, j = 1,2)$ are non negative with $\tau = \max(\tau_{1J})(i, j = 1,2$.as is shown in the appendix, the following system dominates (6.1*)*

$$\left.\begin{aligned} \frac{dZ_1(t)}{dt} &= -m_{11} Z_1(t) + b_{11} Z_1(t-\tau_{11}) - m_{12} Z_2(t-\tau_{12}), \\ \frac{dZ_2(t)}{dt} &= -m_{22} Z_1(t) + m_{21} Z_1(t-\tau_{21}) - b_{22} Z_2(t-\tau_{22}), \end{aligned}\right\} \; t > 0, \quad (6.2)$$

$Z_i(s) = |Y_i(s)|, \qquad i = 1,2 \;\; for \; s \leq 0 \;\; for \; which \; Y_i(s) \; is \; defined.$

Thus the asymptotic stability of the trivial solution of (6.2) implies that of (6.1). it is also shown in the appendix that the solutions of (6.2) are nonnegative for all t $\geq 0$. we will use the following vector-matrix notation:

$$Z = (Z_1, Z_2)^{\tau}, \qquad ||Z|| = |Z_1(t)| + |Z_2(t)| = Z_1(t) + Z_2(t)$$

$$B = \begin{bmatrix} b_{11} & m_{12} \\ m_{21} & b_{12} \end{bmatrix}, \quad ||B|| = max \begin{matrix} b_{11} + & m_{21} \\ b_{21} + & m_{12} \end{matrix}, \tag{6.3}$$

$$A = \begin{bmatrix} -m_{11} & 0 \\ 0 & -m_{22} \end{bmatrix}.$$

The logarithmic norm or the measure of any square matrix B, denoted by $\mu(B)$is defined as follows:

$$\mu(B) = \lim_{\theta \to 0+} \frac{/\!/I + \theta B/\!/ - 1}{\theta}$$

($I$ being an identity matrix of the same size as $B$).The matrix measure compatible with the norms in (6.3) is given by

$\mu(B) = \max_J [b_{jj} + \sum_{\substack{i=1 \\ i \neq J}}^{2} | b_{jJ} | ]$ $B=(b_{1J})$.

We refer to Coppel [11]or Vidyasager [31]for a number of properties of matrix measure We have shown in the appendix that the derivative of $H(Z_1,Z_2)(t)$(seeA.10)satisfies the relation

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq [-m_{11}+b_{11}+m_{21}]Z_1(t)+[-m_{22}+b_{22}+m_{12}]Z_2(t)$, and hence we have

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq \{\mu(A) + /\!/B/\!/\}\{ Z_1(t)+ Z_2(t)\}$

If we assume for the matrix norm $\|B\|$ as in (6.3)that

$$\mu(A) + /\!/B/\!/ < 0 \; or \; -\mu(A) > /\!/B/\!/$$

Then from the results of the appendix (Lemmas 5 and 6) it will follow that $Z_1(t)+Z_2(t) \to 0$ as $t \to \infty$ and thus we have the following:

Lemma 6.1 Assume that the Coefficients of the linear system (6.1)satisfy the condition (6.7). Then the trivial solution of(6.1)is asymptotically stable.

When (6.7)holds an estimate for the decay rate of solutions of (6.1)can be obtained as follows A necessary and sufficient condition for the asymptotic stability of the trivial solution is that all the roots of the characteristic equation

Det $G(\lambda,T_{iJ})=0$

Have negative real parts where

$G(\lambda,T_{1J})=[\begin{matrix} \frac{\lambda+m_{11}}{0} & \frac{0}{\lambda+22} \end{matrix}] - b_{11}^{e-\lambda T_{11}}$ $\begin{matrix} -m_{12}^{e-\lambda T_{12}} \\ -m_{21}^{e-\lambda T_{21}} \quad b_{22}^{e-\lambda T_{22}} \end{matrix}$

$=\lambda I - A - B(\lambda, T_{ij})$(say),

A positive number $\mathfrak{a}$(if it exists)is said to be a decay rate of solution of (6.1) if and only if every root of (6.8)satisfies the condition

$\mathrm{Re}(\lambda)+\mathfrak{a} < 0.$

When (6.11)holds the solution of (6.1)will be of the order of $e^{-at}$ and hence the supremum of all such $\mathfrak{a} > 0$ $a$ satisfying (6.11)will be a desirable parameter to estimate However such a supremum of all such (6.11)may not be possible to determine since it has to be done directly from (6.8)-(6,10)(see also Section 3).We will obtain an estimate of $\mathfrak{a}$ (not necessarily the supremum)I (6.11)by the following procedure.

Replace $\lambda$ byz-$\mathfrak{a}$ in (6.8)and define

$G(z-\mathfrak{a}, T_{IJ}) = M(z) = (z-a)I - A - B(z-\mathfrak{a}, T_{IJ}).$

By Lemma 6.1 a Sufficient condition for all roots of M(z)=0 to have negative real parts equivalently to satisfy

$\mathrm{Re}(\lambda) < -a$

is the following :

$-\mu(aI + A) > \|B_1(a, T_{ij})\|,$

Where

$$B_1(a, T_{ij}) = \begin{pmatrix} b_{11}^{\,e-\lambda T_{11}} & -m_{12}^{\,e-\lambda T_{12}} \\ -m_{21}^{\,e-\lambda T_{21}} & b_{22}^{\,e-\lambda T_{22}} \end{pmatrix}$$

Since $\mu(\mathfrak{a} I + A) = \mathfrak{a} + \mu(A)$ we have from (6.13) that

$\mathfrak{a} + \mu(A) + \|B_I(\mathfrak{a}, T_{IJ})\| < 0$

Is a sufficient condition for Re $(\lambda) < -a$If we denote the left side of (6.15)by $\mathfrak{a} f(\mathfrak{a}, T_{Ij})$ we note that

$f(0, T_{Ij}) = \|B\| + \mu(A) < 0$ $if$ (6.7)$holds$;

Al so $\partial f / \partial \mathfrak{a} > 0, f$ is continuous and monotonic in $(\mathfrak{a})$;$\to \infty as\ \mathfrak{a} \to \infty$; thus by the continuity of $f$ on $\mathfrak{a} \geq 0$, there exists an interval of the from $[0, ^{*})$such that

$f(0, T_{Ij}) < 0$ $for\ \mathfrak{a} \in [0, \mathfrak{a}^{*})$

And hence

$f(\mathfrak{a}, T_{Ij}) = 0$

Thus an upper bound $\mathfrak{a}^*$ of $\mathfrak{a}$ for which we have $f(\mathfrak{a}, T_{Ij}) < 0$ $\mathfrak{a} \in [0, \mathfrak{a}^*)$ can be obtained as the unique (by the monotonicity of $f$ on $\mathfrak{a}$) positive root of$\max \left\{ b_{11}^{e-\lambda T}{}_{11} - m_{12}^{e-\lambda T}{}_{21} + \mathfrak{a}\text{-}\min \left\{ m_{11} \right.\right.$

$\left.\left. b_{22}^{e-\lambda T}{}_{22} - m_{21}^{e-\lambda T}{}_{12} \right\} \right\} m_{22}$

Where the maximum in (6.18) is taken over $b_{11}, b_{22}, m_{12}, m_{21}$ and $T_{ij}\,(i, = j\ 1,2)$. A numerical procedure can be devised to solve equations of the form (6.18) in concrete cases. If one is interested in a smaller bound on the decay rate it can be obtained as the unique root of

$$\| B \| e^{ar} + a + \mu(A) = 0, \tau = \max_{i,J} T_{ij}. \qquad (6.19)$$

It will follow from (6.19) than an increase in T

Will reduce the value of the root of (6.19). We remark that since our estimate $a^*$ of dacay rate is based on a sufficient condition of the form (6.7), we can conclude that the solotions of (6.1) will converge at least as fast as $e^{-a*t}$ where $a^*$ is the unique root of either of (6.18) or (6.19).

## Implications

Authors have formulated a class of two systems model systems staring from the elementary principal that the rate of change of a population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) equals the creation rate minus the destruction rate. This is a departure from the conventional model building in population ecology where one usually starts from the Malthusian formulation, now known as the Kolmogorov formulation in the nonlinear context. In our analysis of the models proposed we have obtained sufficient conditions for the global existence of nonnegative solutions of the model systems (this is trivial in the Kolmogorov formulation) and the other desirable properties such as the existence of suitable steady states. The two classes of models proposed for competition and debt contain intraspecific self – regulation biologists that each systems is under the action of a self-regulating negative feedback force, opinions and arguments differ regarding the mechanism by which such an effect is realized; we refer for more details on this aspect to a collection of articles compiled by Tamarin [30].

Gopalasamy's main result is the following : if a self –regulating negative feedback mechanism (intraspecific oppugn ant) can act without time delay and if such an effect can keep an ecosystem ( a community of one or more species of populations) asymptotically stable ( at least locally about a non trivial steady state, then time delays of any magnitude in any other interaction involving positive or negative feedbacks cannot destabilize such a stable ecosystem ( in a mathematical sense); that is such delays cannot cause a hopf-type bifurcation leading to persistent and undamped fluctuations. The delay-induced bifurcations to oscillations in models have been investigated by many authors. (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Qualitatively result is not new to the literature the effects of delays on otherwise stable and unstable systems have been considered by Cushing[12]where (pages 35-39)he has shown that an otherwise unstable system can be made

less unstable by delays, Such an effect of delays on unstable systems appear to have been first proposed by Beddinton and May [2]who have also shows that delays can make stable systems less stable leading to a conclusion "time delays are not necessarily destabilizing" It is also shown by Cushing([12], pages 75-103)that while time delays cannot stabilize(in the mathematical sense)an otherwise unstable two species system in competition or mutualism time delays in there species systems can stabilize otherwise unstable systems (even in a mathematical sense)Cushing and Saleem [14]have formulated a model where delays in the positive feedback can have a stabilizing effect: Hale[18]discusses briefly the stabilizing nature of delays and given an example to show that an increase in delays can make an unstable system stable Thus for a better understanding of the effects of delays on otherwise stable systems one has to estimate the rate of decay of perturbations (or the characteristic return time )in terms of the parameters of the system especially when the delays do not cause a mathematical instability This aspect has been investigated by a procedure slightly different from that of the characteristic equation we have proposed (section6) method for the estimation of decay rate of perturbations when there are multiple delays Since many species systems with multiple delays cannot be analysed (due to complexity)by the characteristic equation our methods using comparison lead to easily verifiable sufficient conditions The sufficient conditions obtained are in the form of a diagonal dominance type; these conditions can be weakened to quasi-diagonal dominance type as in Ladde [20], Siljak[20]Anderson [1], Lewis and Anderson [21,22](Ladde's result [20] is delay dependent due to the fact that delay appears in self –regulation also).Verification of quasi –diagonal dominance conditions in arbitrary matrices (except for Metzler matrices )is an impossible task and quasi –diagonal dominance is hard to interpret.

If follows from our discussion that those confederation systems which display delay induced instability leading to permanent and undamped oscillations should have one or more of the following aspects:

(1)time delays are in intraspecific interaction terms (and possibly in other terms also): (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

(2) Intraspecific self-regulation terms are absent in density dependent from in one or more populations of the system;

(3) Intraspecific negative feedback effect are not strong enough to maintain the stability of the relevant steady state of the nondelayed system.

We can now infer how the delay induced stability leads to fluctuations; a primary mechanism of such fluctuations can be described at least theoretically as follows : in its attraction of a system to a steady state, intraspecific self-regulation fails to be effective and the system overhoots the steady state and at such higher densities intraspecific competition becomes operationally effective leading the system to understood the steady state and this process is repeated again and again resulting in undamped oscillations about the steady state. It follows from our results that if intraspecific competition is sufficiently strong, the system cannot become oscillatory even when there are time delays, so long as they do not interfere with the strength of the intraspecific self – regulation.

We are thus led to the question of existence and importance of intraspecific competition in population systems. In the introduction to an anthology of articles by several authors, Tamarin [ 30] begins with the statement "all animal populations have the reproductive capacity to increase without limit, yet none di" and in his concluding remarks writes " this leaves us with only one great concept in the area of confederation do main regulation: that all credit-debit are regulated" More specifically this is applicable to all the confederation of (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions,

financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

In his discussion on the effect of competition on the densities of populations, Nicholson ([ 25], $page$ [ 147] ) has reached the conclution "Intraspecific competition automatically regulates the severity of its action to the requirements of each denizen , provided the inherent resistance of the environment is sufficiently low to permit the species to exist." In order to look for a mechanism by which intraspecific competition plays its role of decreasing growth rate as confederation density increases, Calhoun [ 7] set up an experiment with Norway rates (Rattus norvegicus) and found that increased density leads to increased hostility which disturbs the psychology and physiology causing decreased reproductive output. Many other mechanisms by which intraspecific compition self-regulate populations have been put forward by a number of authors (see Tamarin [30]). This is true of (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

If we accept the view that intraspecific competition is really predominant in population regulation, then we can understand how and why time delays destabilize stable system leading to oscillations; time delays render intraspecific competitive regulation territorially (in case of particles we take the field in which it operates) weak so that self –regulatory forces become temporarily insufficient. If we now agree that all population models should incorporate self-regulating intraspecific competition terms in surfeit strength along with other tropic intraspecific competition terms in sufficient strength along with other atrophic interactions then it is not fair to say " time delays destabilize" and we note from our analysis that it is usually a failure in self- regulation which can cause instability rather than delays by themselves.

We will give one example to illustrate our point. Consider a credit-debit system in which the debits is self –regulating by a logistic load as follows: (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

$$\frac{dx}{dt} = r\left(x - \frac{x^2}{K}\right) - axy = b_I(x) - m_I(x,y) \quad (say),$$

$$\frac{dy}{dt} = \beta xy - \delta y = b_2(x,y) - m_2(y) - \quad (say),$$

Where $K, a, \beta, \delta$ are positive constants such that the steady state $x^*\ y^*$ exists where

$$x^* = \frac{\delta}{\beta}; \quad y^* = \frac{r}{a}\left(-\frac{\delta}{K\beta}\right) > 0.$$

In this example the predator self –regulation $(\partial m_2/\partial y)$ is not sufficiently strong in the sense

$$x^* = \left(\frac{\partial m_2}{\partial y}\right)_{(x^* y^*)} \succ \left(\frac{\partial b_2}{\partial y} + \frac{\partial m_1}{\partial y}\right)_{(x^* y^*)}$$

If in such a model there is a time delay in the credit creation rate, such a delay can lead to instability of $(x^*\ y^*)$ of setting up undamped oscillations. If however we modify the growth rate by the inclusion of a density dependent self-regulation so that credit

$$\frac{dy}{dt} = \beta\, xy - \delta y - \eta y^2$$

For some positive constant η then it is possible, depending on η, for such a modified system to remain stable without oscillations even if there is a time delay in the credit creations rate. We have another completely different justification to have $\eta y^2$ in the credit dynamics: for instance if the debit is removed from this cannot happen. While this example serves to illustrate an aspect of our result, our advocacy of the significance is not based on this example.

Gopalasamy has thus derived conclusion that "time delays need not destabilize" in models with several discrete delays. However we have provided in the appendix the necessary tools (Lemmas 3,4,5) to prove all our results when the relevant delays are all (or some) continuously distributed. Some authors (Caswell [9], May [23], Cushing [12] appear to regard continuously distributed dalays as financially/cosmologically credit wise obtained when he subjected algae to a variable nitrate environment and obtained a better fit for his data with a continuous delay. The questions of whether a discrete or continuous delay is more realistic may well depend on what process is being modeled. For instance an increased supply of some resource at some point in time may trigger reproduction at a particular moment in time later irrespective of the conditions of the resource during the intervening period. If the duration of increased resource supply necessary to elicit such a response is short compared to the intervening period before the reproduction occurs, then a discrete time delay may be appropriate. In a large number of cases of animal. Note fungus algae are prey. Here it shall be an account which loses the energy, amount, or matter for that matter. Same result holds good for (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Populations of particles, atoms, molecules, propositions, systems types chosen the past dietary and nutritional history of an animal over a long period plays an influential role in determining the current bahaviour of the dietary and nutritional history of an animal over a long period plays an influential role in determining the current behavior of the animal and in such cases continuous time delays will be appropriate.

In conclusion the following comments are appropriate. The time delays since the current creation rate is usually due to a contribution of those mature animals living a few units of time earlier depending on the type of animal or organism; the time delays appearing in interspecific interactions can be due to the fact that in order to successfully interact and compete interferentially a developmental maturity may be necessary. In the case of predation, delays in the birth rate of prey can be due to delayed maturation while that in the birth rate of a predator is due to gestation delay with delayed fecundity; the delays in the predation rate are again are again as before due to the maturity needed for a predator to successfully organize and succeed in an attack.

In our analysis we have assumed throughout a temporally uniform environment neglecting dispersal of populations and their age structure. We expect our results to hold under suitable conditions both in temporally periodic environments with the steady state of a constant environment being replaced by a periodic solution

(nonstationary steady state) and when the effects of dispersion and age structures are considered. It is worthwhile to modify our purely ecological, cosmological, financial, propositional, problematic systems models so as to provide them with a genetic foundation and examine the effects of time delays since forces of selection act with some time delay. Such a generalization will definitely lead to a better understanding of "how confederations regulate themselves".

## Appendix

First we derive a comparison result.

Lemma 1. Consider the delay differential system

$$\left.\begin{aligned}\frac{dY_1(t)}{dt} &= -m_{11}\ Y_1(t) + b_{11}Y_1(t-\tau_{11}) - m_{12}\ Y_2(t-\tau_{12}),\\ \frac{dY_2(t)}{dt} &= -m_{22}Y_2(t) - m_{21}Y_1(t-\tau_{21}) + b_{22}\ Y_2(t-\tau_{22}),\end{aligned}\right\} (A.1),$$

Where the $m_{ij}$, $b_{ij}$ $(i.j = 1,2)$ are positive while $\tau_{ij}$ $(i.j = 1,2)$ nonnegative with $\tau = \max \tau_{ij}$, $i,j = 1,2$. Then

$$\left.\begin{aligned}\frac{d|Y_1(t)|}{dt} &\le -m_{11}\,|Y_1(t)| + b_{11}|Y_1(t-\tau_{11})| + m_{12}\ Y_2(t-\tau_{12})|,\\ \frac{d|Y_1(t)|}{dt} &\le -m_{22}\,|Y_2(t)| + m_{21}|Y_1(t-\tau_{21})| + b_{22}\ |Y_2(t-\tau_{21})|,\end{aligned}\right\} (A.2)$$

Where

$$\frac{d}{dt}\,|\,Y_1(t)\,| = \left\{\begin{aligned}&\frac{dY_1(t)}{dt}\ if\ \ Y_1(t) \ge 0,\\ &\frac{dY_1(t)}{dt}\ if\ \ Y_1(t) < 0,\end{aligned}\right\}\ i = 1,2. \qquad (A.3)$$

The proof is by direct verification.

LEMMA 2. Consider now (A.2) and the following system

$$\left.\begin{aligned}\frac{dZ_1(t)}{dt} &= -m_{11}\ Z_1(t) + b_{11}Z_1(t-\tau_{11}) + m_{12}\ Z_2(t-\tau_{12}),\\ \frac{dZ_2(t)}{dt} &= -m_{22}Z_2(t) + m_{21}Z_1(t-\tau_{21}) + b_{22}\ Z_2(t-\tau_{22}),\end{aligned}\right\} \qquad (A.4)$$

Together with the initial conditions

$$Z_1(s) > |Y_1(s)|,\ \ s \in [-\tau, 0],\ i = 1,2. \qquad (A.5)$$

Then

$$|Y_1(t)| < Z_1(t) \text{ for } t \ge 0 \quad \text{and} \quad i = 1,2. \qquad (A.6)$$

PROOF. Let $U_1(t) = Z_i(t) - |Y_1(s)|$, $i = 1,2.$ $t \geq -\tau.$ Then we have from (A.2) and (A.4),

$$\left.\begin{aligned} \frac{dU_1(t)}{dt} &\geq -m_{11}\ U_1(t) + b_{11}U_1(t-\tau_{11}) + m_{12}\ U_2(t-\tau_{12}), \\ \frac{dU_2(t)}{dt} &\geq -m_{22}U_2(t) + m_{21}U_1(t-\tau_{21}) + b_{22}\ U_2(t-\tau_{22}), \end{aligned}\right\} t > 0, \quad (A.7)$$

With

$$U_1(s) = Z_i(s) - |Y_i(s)| > 0, \quad i = 1,2, s \in [-\tau, 0]. \quad (A.8)$$

Consider (A.7)-(A.8) first on the interval $[\tau, 0]$ : It will follow from (A.7) that $U_1(t) \geq U_1(0) \exp [-m_{11}t]$

$$+\int_0^t [b_{11}U_1(s-\tau_{11}) + m_{12}\ U_2(s-\tau_{12})] \exp [-m_{11}(t-s)ds > 0,] \quad (A.9)$$

For all $t \in [0, \tau^*]$ since $U_i(s) > 0$ on $[-\tau^* 0]$ where $\tau^*$ = positive minimum of $(\tau_{11,}\tau_{12,}\tau_{21,}\tau_{22,})$. Similarly $U_2(t) > 0$ for $t \in [0, \tau^*]$. Now by a similar argument it will follow that $U_1(t) > 0, U_2(t) > 0$ for $t \in [\tau^*, 2\tau^*]$. Repeating the above procedure we obtain the result (A.6).

It follows from Lemmas 1 and 2 that in order to prove theorem 3.1 it is enough to prove that $Z_i(t) \to 0$ as $t \to \infty$ for $i = 1,2$ . which we will do as follows.

PROOF OF THEOREM 3.1 Define a Lyapunov-like function

$$H(Z_1,Z_2)(t)=Z_1(t)+Z_2(t)$$

$$+\int_{1-T}^t b_{11}\, Z_1(s)ds + \int_{1-T}^t m_{12}\, Z_2(s)ds$$

$$+\int_{1-T}^t m_{21}\, Z_1(s)ds + \int_{1-T}^t b_{22}\, Z_{22}(s)ds$$

Since each term of H in (A.10)is nonnegative for t>0, if we show that $H \to 0$ as $t \to \infty$ it will follow from Lemma 2 that

$$0 \leq |\, Y_{1,}(t)\,| \leq Z_1(t) \to 0 \text{ as } t \to \infty,$$

Implying that the perturbations $Y_{1,}$, $Y_2$, governed by (A.1)approach zero as $t \to \infty$, and this will prove our result Calculate the rate of change of H along the solutions of (A.4)so that

$$\frac{dH}{dt} = \frac{\partial H}{\partial t} + \frac{\partial H}{\partial Z_1}\frac{dZ_1}{dt} + \frac{\partial H}{\partial Z_2}\frac{dZ_2}{dt}$$

$$= -[Z_1(t)\{m_{11} - (b_{11} + m_{12})\} + Z_2(t)\{m_{22} - (b_{22}m_{12})\}] < 0$$

On using (2.10)since $Z_1(t),Z_2(t)$ are positive for all t>0. Thus $H(Z_1, Z_2)(t) \to 0$ as $t \to \infty$ since H is bounded below by zero: it follows from Lemma 6 below that $Z_1(t) \to 0, Z_2(t) \to 0$ as $t \to \infty$ the proof is complete.

The following set of results are intended to be a guide to prove result similar to the above in the case of continuously distributed delays An analogue of Lemma I is the following.

LEMMA 3 Let $K_{ij}(s)$,s∈[0,∞), $i,j = 1,2$,denote nonnegative continuous functions integrable on [0,∞) normalised such that

$\int_0^\infty sk_{ij}(s)ds< \infty$; $\int_0^\infty k_{ij}(s)ds=1$, I,j=1,2.

Consider the intergrodifferential system

$$\frac{dY_1(t)}{dt} = -m_{11}Y_1(t) + b_{11}\int_0^\infty k_{ij}(s)\, Y_1(t-s)ds$$

$$-m_{12}\int_0^\infty k_{12}(s)\, Y_2(t-s)ds$$

$$\frac{dY_2(t)}{dt} = -m_{22}Y_2(t) + m_{21}\int_0^\infty k_{21}(s)\, Y_1(t-s)ds$$

$$+b_{22}\int_0^\infty k_{22}(s)\, Y_2(t-s)ds$$

*Where* $m_{ij}, b_{ij}$ $(i,j = 1,2)$ *are positive constants Then*

$$\frac{d|Y_1(t)|}{dt} \leq -m_{11}|\,Y_1(t)\,| + b_{11}\int_0^\infty k_{11}(s)|Y_1(t-s)|ds,$$

$$+m_{12}\int_0^\infty k_{12}(s)|\,Y_2(t-s)|ds,$$

$$\frac{d|Y_2(t)|}{dt} \leq -m_{22}|\,Y_2(t)\,| + m_{21}\int_0^\infty k_{21}(s)|Y_1(t-s)|ds,$$

$$+b_{22}\int_0^\infty k_{22}(s)|\,Y_2(t-s)|ds,$$

*Where* $d|Y_1(t)||/dt$ *is defined by* (A.3).

Proof is again by direct verification.

The proof of the following lemma is similar to that of Lemma 2 above.

Lemma 4. *Consider (A.14) and the following intergrodifferential system:*

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1(t) + b_{11}\int_0^\infty k_{11}(s)Z_1(t-s)ds$$

$$+m_{12}\int_0^\infty k_{12}(s)Z_2(t-s)ds,$$

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1(t) + b_{11}\int_0^\infty k_{11}(s)Z_1(t-s)ds$$

$$+b_{22}\int_0^\infty k_{22}(s)Z_2(t-s)ds,$$

Together with the initial conditions

$$\infty > Z_i(s) > |Y_i(s)| for\ s \in (-\infty, 0).$$

(A.16)

Then $|Y_i(t)| < Z_i(t) for\ t \geq 0, i = 1,2.$

The following lemma will be useful to prove a result analogous to that of Theorem 3.1 corresponding to a class of continuously distributed time delays.

LEMMA 5. Let $Z_1(t), Z_2(t)$ be the solutions of (A.15)- (A.16) where $(Y_1\ Y_2)$ corresponds to an arbitrary but fixed solutions of (A.13). Then as a consequence of Lemma & it will follow that

$$0 \leq |Y_i(t)| \ Z_i(t) \to \infty \ as\ t \to \infty, \qquad i = 1,2,$$

(A.17)

Whenever

$$\begin{cases} M_{11} > b_{11} + m_{21}, \\ M_{22} > b_{22} + m_{12} \end{cases}$$

PROOF. Define a Lyapunov –like function H where

$$H(Z_1,Z_2)(t) = Z_1(t) + Z_2(t) + b_{11}\int_0^\infty k_{11}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$$

$$+ m_{12}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$$

$$+ m_{21}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$$

$$+ b_{22}\int_0^\infty k_{22}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$$

Where $(Z_1,Z_2)$is any solution of (A.15)-(A.16) for arbitrary $(Y_1,Y_2)$defined on $(-\infty, 0]$ Computing the rate of change of H in (A.19)along the solution of (A.15)we derive that (on using (A.12))

$$\left(\frac{dH(Z_1Z_2(t)}{dt} = [Z_1(t)\{m_{11} - (b_{11} + m_{21})\} + Z_2(t)\{m_{22} - (b_{22} + m_{12})\}\right]$$

$< 0$(using (A.18)) for $Z_1(t) > 0\ and\ Z_2(t) > 0.$

By Lemma 6 below the result will follow

Lemma 6. Let H(t)=H$(Z_1, Z_2)(t)$ be as defined in (A.10) Suppose

$$m_{11} > b_{11} + m_{21} \text{ and } m_{22} > b_{22} |m_{12}$$

Then

$$Z_1(t) + Z_2(t) \to 0 \text{ast} \to \infty.$$

PROOF. Suppose (A.21) dose not hold; then there exists at least one sequence

$\{t_n\} \to \infty$ as n $\to \infty$ such that

$$[m_{11} - (b_{11} + m_{21})]\, Z_1(t_n) + [m_{22} - (b_{22} + m_{12})]Z_2(t_n) > \varepsilon \qquad (A.22)$$

For some sufficiently small positive number $\varepsilon$ and hence

$$\frac{dH(t_n)}{dt} < -\varepsilon, n = 1,2,3, \ldots.$$

By the continuity and differentiability of H with respect to t,

$$\mathrm{H}(t_n - \varepsilon) - \mathrm{H}(t_n) = \frac{dH(t^*)}{dt}(-\varepsilon), \quad t_n - \varepsilon < t^* < t_n$$

If $\varepsilon$ is small we will have

$$= \frac{dH(t^*)}{dt}\varepsilon/2$$

So that

$$\mathrm{H}(t_n) \leq \mathrm{H}(t_n - \varepsilon) - \varepsilon^2/2$$

$$\leq H(t_{n-1}) - \varepsilon^2/2$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2)$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2):$$

$$\leq H(t_0) - \mathrm{n}(\varepsilon^2/2): \mathrm{t}_0 = 0.$$

(A.12) and Lemma 5 will imply that $H(t_0) < \infty$ for any $t_0 > 0$. If n is sufficiently large it will follow that $H(t_n) < 0$ since $H(t_0) < \infty$, which contradicts the nonnegativity of H for all $t \geq 0$. Hence $Z_1(t) + Z_2(t) \to 0$ as $t \to \infty$ implying the asymptotic stability of the trivial solution of (A.1)(since $Z_1$ dominates $y_1$)

We note the result of Lemma 5 and 6 will remain valid if (A.18) is replaced by the condition $\mu(A) + \|B\| < 0$

Where A and B are defined as in (6.3).

We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities,

apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems

**Gopalsamy**

# Model For Innocuously Benign Delay In A Periodic System

**Abstract**

Sufficient conditions are obtained for the existence of a unique linearly denomination stable positive periodic solution of a system model of two competitions in a periodic environment with time lags in interspecific interactions. It is shown that if the self- regulating intraspecific interaction effects are strong enough and act without time delays then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a periodic domination. An example in finance will make things clear. Corporate finance for one big plant generates lot of accounts of the workers whose salary shall be disbursed to the debit of cash credit account. Here Cash credit account is a resource as also the one to which debits are raised, a la predator. Some of the accounts from which the credits come are demand bills purchase, usance bills discount, vostro and nostro accounts in respect of foreign bills and Letters of credit. In economics inventories are used to produce output, output is sold to produce income which might be saved in part and recycled in investment. Cosmologically take an example of a Black hole which gormandizes matter and radiates energy. Notwithstanding information paradox, one thing is clear that matter is being created in to radiation. We show how the model could be applied to axiomatic predications, postulation alcovishness, phenomenological methodologies, propositional statement, and other problems in cosmology, philosophy, and even in Green House effects. We study the sine qua non properties of such a system. In fact these models can be applied to such unsolved problems like Yang mills Theory and mass gap, and other problems. As we apply the model, we shall go on dilating the position status and the applicability of the model. We shall not touch the electrical engineering problems which might rise questions like what is charge of a Bank. We may say when a customer enters a Bank; we can say the Bank is charged. These factors are still relativistic and we shall make an in-depth study as we progress. Gopalsamy's research was published in Australian Journal of Mathematical Society vide: B 25349=365 It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 7. Introduction

At the outset we assume that the system satisfies Haimovici is caution. The purpose of this article is to establish a set of sufficient conditions for the existence of a unique linearly asymptotically stable positive periodic solution for a nonautonomous system of the form

$$\frac{dx(t)}{dt} = x(t)f_1\big(t, x(t), y(t-\tau_1)\big),$$

**(1.1)**

$$\frac{dy(t)}{dt} = y(t)f_2\,\big(t, x(t-\tau_2), y(t)\big).$$

In the context of confederation dynamics a model system of the form (1.1) means that the interspecific interactions occur with time delays $\tau_1$ 0, $\tau_2$ 0 in a periodic environment, if $f_1$ *and* $f_2$ are assumed to be periodic in $t$ with a common period $\omega$ 0. Note the interactions are inter specific.

Although model systems in periodic domains have been investigated by a large number of authors (for some references see Gopalsamy (4], [5]), periodically varying environments in time delayed model systems of two or more interacting species have not been investigated. The need for Multitype system models in periodic environments and time delayed interactions are evident from the extensive literature of that type relating to single species. While it is known that periodic ecosystems can support more species than uniform environments (Gopalsamy [4]), time delays might destabilize the system leading to a possible loss of one or more species. Furthermore, periodic environments in population dynamics have some relevance to natural selection and evolution; we will consider this aspect briefly at the end.

The principal finding of the following analysis of (1.1) is that, under suitable hypotheses on $f_1$ *and* $f_2$, the system (1.1) has a linearly asymptotically stable unique positive periodic solution for all nonnegative delays $\tau_1, \tau_2$ (while our numerical experiments suggest that the periodic solution of (1.1) is globally asymptotically stable and is independent of the delay parameters $\tau_1, \tau_2$ it is not known how one can prove such a strong result). Since time delays are usually associated with possible destabilization of otherwise stable systems, we can describe the time delays appearing on (1.1) as “harmless”.

It has been recently shown by the author (Gopalsamy [6]) that if the self –regulating negative feedback in each denomination is strong enough and acts without time delays, then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a temporally uniform domain. The result of this paper is thus an extension to a periodic environment of the "harmless" nature of interspecific interaction delays in a temporally constant environment mentioned above.

## 8. Existence and uniqueness of periodic solutions

We will now formulate the problem of existence of periodic solutions of the periodic delay differential system

$$\frac{dx_1(t)}{dt} = x_1(t) f_1\big(t, x_1(t), x_2(t-\tau_{12})\big),$$

$$t > 0,$$

**(2.1)**

$$\frac{dx_2(t)}{dt} = x_2(t) f_2\big(t, x_1(t-\tau_{21}), x_2(t)\big),$$

(Where $\tau_{12}, \tau_{21}$ are nonnegative constants) in a form suitable for the application of fixed point theory of positive operators. We make the following assumptions regarding $f_1$ and $f_2$ in (2.1).

($H_1$) The partial (Frechet) derivatives of $f_1, f_2$ satisfy the following:

$$\frac{\partial f_i}{\partial x_j} < 0, \qquad i, j = 1, 2 \text{ For all } t > -\tau \text{ where } \tau = \max(\tau_{12}, \tau_{21}). \qquad \textbf{(2.2)}$$

($H_2$) $f_1(t, x_1, x_2) \equiv f_1(t + \omega, x_1, x_2)$ for all $t > -\tau, x_1 \geq 0, x_2 \geq 0$. **(2.3)**

($H_3$) There exist positive constants $b_{il}$, $b_i^u$, $a_{il}$, $a_i^u$ ($l = 1, 2$) such that

$$b_{il} - a_{il} x_i < f_i(t, x_1, x_2) < b_i^u - a_i^u x_i,$$

$$b_{il}/a_{il} < b_i^u/a_i^u, \qquad \textbf{(2.4)}$$

For $i$= 1, 2 and $(x_1, x_2) \in R_+ \times R_+$. If $\tau = 0$ then one can show that (2.1)-(2.4) will imply that the nondelayed periodic system (2.1) will have an asymptotically stable strictly positive periodic solution (for details see Gopalsamy [5]); this result will become a special case of what we obtain in the following.

We will use the following notation:

$$B_\omega^\tau = \{h: [-\tau, \infty) \to R, h \text{ is continuous}, h(t + \omega) = h(t), t \geq -\tau\},$$

$$B = B_\omega^\tau \times B_\omega^\tau,$$

$$K_1 = \{ h \in B_\omega^\tau | h(t) \geq 0, t \in [-\tau, \infty)\},$$

$$K = K_1 \times K_1,$$

$$||h|| = \left\{ \sup_{t \in [0,\omega],} | h(t)|, h \in B_\omega^\tau \right\}$$

$$||x_1, x_2|| = \max\{||x_1||, ||x_2||, (x_1, x_2) \in B\},$$

$$G_1 = \{(x_1, x_2) \in K, ||(x_1, x_2)|| \leq b_l = \min(b_{1l} / a_{1l}, b_{2l}/a_{2l})\}$$

$$G_2 = \left\{(x_1, x_2) \in K, ||(x_1, x_2)|| \geq b^u = \max\left(\frac{b_1^u}{a_1^u}, \frac{b_2^u}{a_2^u}\right)\right\},$$

$\bar{G} =$ closure of $G_2$ *in* $B$ and

$\frac{\bar{G}_2}{G_1} =$ complement of $G_1 in\ \bar{G}_2$; $G_1 \subset G_2$ by **(2.4)**

The cone $K = K_1 \times K_1$ induces a partial order in $B$ and we will write $(x_1, x_2) \leq (y_1, y_2)$ if and only if $x_1 - y_1 \in K_1$. we define an operator $N$ on $B$ as follows

$$N(x_1, x_2)(t) = \begin{matrix} N_1(t, x_1, x_2) \\ N_2(t, x_1, x_2) \end{matrix} \quad \text{for t> } 0, \qquad \textbf{(2.5)}$$

Where

$$N_1(t, x_1, x_2) = \int_0^\omega G_{11}(m, \omega, t-s) g_1\left(m, s, x_1(x) x_2(s - \tau_{12})\right) ds$$

$$+ \int_0^1 G_{12}(m, 0, t-s) g_1\left(m, s, x_1(s), x_2(s - \tau_{12})\right) ds,$$

$$N_2(t, x_1, x_2) = \int_0^\omega G_{22}(m, \omega, t-s) g_2\left(m, s, x_1(s - \tau_{21}), x_2(s)\right) ds \qquad \textbf{(2.6)}$$

$$+ \int_0^1 G_{21}(m, 0, t-s) g_2\left(m, s, x_1(s - \tau_{21}), x_2(s)\right) ds,$$

For $t > 0$, with

$$g_1\left(m, s, x_1(s), x_2(s - \tau_{12})\right) = m x_1(s) + x_1(s) f_1\left(s, x_1(s), x_2(s - \tau_{12})\right),$$

**(2.7)**

$$g_2\big(m,s,x_1(s-\tau_{21}),x_2(s)\big) = mx_2(s) + x_2(s) f_2\big(s,x_1(s-\tau_{21})x_2,(s)\big),$$

$$G_{ll}(m,\omega,t-s) = e^{-m\omega}(1-e^{-m\omega})^{-1}e^{-m(t-s)}, \qquad i=1,2,$$

**(2.8)**

$$G_{ij}(m,0,t-s) = e^{-m(t-s)}, i \neq j, i,j = 1,2,$$

$$N(x_1,x_2)(t) = \begin{cases} N_1(n\omega+t,x_1,x_2) \\ N_2(n\omega+t,x_1,x_2) \end{cases} for\ t \in [-\tau,0], \qquad \textbf{(2.9)}$$

*m* being a suitably selected positive constant (more will be said below) while *n* is a positive constant such that $n\omega > \tau$.

We note immediately the following from the definition of $N$: $if\ (x_1,x_2) \in \bar{G}/G_1$ then $N(x_1,x_2)(t)$ and if $N(x_1,x_2)$ is periodic $\omega$ for $t$ then $N(x_1,x_2)$ is periodic in *t* with period $\omega$ for $t \geq -\tau$. The existence of at least one positive periodic solution of (2.1)-(2.4) is established by means of a few lemmas and a result of Schmitt [9].

Lemma 1. *A necessary and sufficient condition for* $(x_1(t),x_2(t))$, $t \geq -\tau$ *to be a periodic solution of (2.1)-(2.4) of period* $\omega$ *is that* $(x_1,x_2)$ *be a fixed point of N such that*

$$N(x_1,x_2)(t) = \big(x_1(t),x_2(t)\big) \quad and\ (x_1,x_2) \in K \cap \left(\frac{\bar{G}}{G_1}\right).$$

PROOF. The details of proof are similar to those of Lemma 2 of Busenberg and Cooke [1]. We first verify the continuity of $N(x_1,x_2)$ assuming that $N(x_1,x_2)(t)$ is periodic in *t* of period $\omega$. The continuity of $N(x_1,x_2)$ for $t > 0$ follows from the definition in (2.5)-(2.8) WHICH ALSO PROVIDES $\lim_{t\to 0+} N(x_1,x_2)(t)$. It will follow from (2.9) and the assumed periodicity of $N(x_1,x_2)$ $that\ N(x_1,x_2)$ is continuous

For $t \geq -\tau$. Now we will verify the periodicity of $N(x_1,x_2)$ in $t$. A simple computation leads to

$$N_1(t+w,x_1,x_2) - N_1(t,x_1,x_2) = e^{-mt}F_1(t),$$

**(2.10)**

Where

$$F_1(t) = \int_0^{\omega} G_{11}\,(m, \omega, t-s\,)g_1\big(m, s, x_1(s), x_2(s-\tau_{12})\big)ds$$

$$+\int_0^{\omega} G_{12}(m, 0, t-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$+\int_{\omega}^{\omega+t} G_{12}(m, 0, \omega-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$-\int_0^{\omega} G_{11}(m, \omega-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$+\int_0^{t} G_{12}(m, 0, -s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

**(2.11)**

Which on simplification using the periodicity of $f_1$ shows that $F_1(t) \equiv 0$ whenever $x_1, x_2, \in B$, thus implying the periodicity of $N_1$ in $t$ of period$\omega$. A similar argument for $N_2$ proves that $(x_1, x_2) \in B \Rightarrow N\,(x_1, x_2) \in B$.

Now suppose $(x_1, x_2)$ is a fixed point of *N and let* $(x_1, x_2)\ \in K \cap (\bar{G}/G_1\,)$. Then it is easily seen that $(x_1, x_2)$ is a solution of (2.1) and since $(x_1, x_2)\ \in K \cap (\bar{G}/G_1\,)$ it will follow that $(x_1, x_2)$ is a periodic solution of (2.1).

Conversely, let $(x_1, x_2)$ be a periodic solution of period $\omega$ of (2.1). Then $(x_1, x_2)$ is necessarily a periodic solution of

$$\frac{dx_1(t)}{dt} + mx_1(t) = g_1\,\big(m, t, x_1\,(t), x_2(t-\tau_{12})\big)$$

**(2.12)**

$$\frac{dx_2(t)}{dt} + mx_2(t) = g_2\,\big(m, t, x_1\,(t-\tau_{21}), x_2(t)\big),$$

For any arbitrary real constant *m* and hence for any fixed constant $m > 0$. It follows from (2.12) that any periodic solution of (2.1) and hence of (2.12) is a fixed point of $N(x_1, x_2)$. We claim that such a fixed point is in $K \cap (\bar{G}_2/G_1)$. Suppose $(x_1, x_2) \notin K \cap (\bar{G}_2/G_1)$. the following possibilities can arise:

(vii) Either $x_1, x_2$ or both $x_1$ *and* $x_2$ may become negative for some value of *t* and it will follow from (2.1) that subsequent value of $x_1, x_2$ or $x_1$ *and* $x_2$ will be negative and hence from (2.4) and (2.1) their derivatives will be negative for all subsequent *t* , implying that $x_1, x_2$ or $x_1$ *and* $x_2$ cannot be periodic.

(viii) While $(x_1, x_2) \in K \subset B$ there may be a value of *t*, say $t^*$, such that $x_1(t^*) > b_1^u/a_1^u$ or $x_2(t^*) > b_1^u/a_1^u$ and in such a case $x_1(t) or\ x_2(t)$ cannot be periodic since respectively $dx_1(t^*)/dt < 0\ or\ dx_2(t^*)/dt < 0,$ which will imply that the values of $x_1(t^*)\ or\ x_2(t^*)$ cannot be repeated from below.

(ix) Similarly if there is a $t^*$ such that $0 < x_1(t^*) < b_{1l}/a_{1l}\ or\ 0 < x_2(t^*) < b_{2l}/a_{2l}$ then such solutions cannot be periodic since $dx_1(t^*)dt > 0$ or $dx_2(t^*)dt > 0$, which will mean that the values of $x_1(t^*)$ or $x_2(t^*)$ cannot be repeted from above.

Thus every periodic solution of (2.1)-(2.4) is a fixed point of *N* in $K \cap \left(\frac{G_2}{G_1}\right) \subset B$.

**LEMMA 2.** *Let N be as in (2.5)-(2.6). Suppose* $\lambda \in (0,1) and\ \lambda N(x_1, x_2) = (x_1, x_2)$ *for some* $((x_1, x_2) \in K - \{0\}$. *then there exist* $\alpha_1^* > 0, \alpha_2^* > 0$ *such that*

$$0 < x_1(t) \leq \alpha_1^* < b_1^u / a_1^u,$$

$$0 < x_2(t) \leq \alpha_2^* < b_2^u / a_2^u,$$

*for* $t \in [-\tau, \infty)$.

PROOF. If $\lambda N(x_1, x_2) = (x_1, x_2)$ then from the definition of N we get for all $t > 0$ and $(x_1, x_2) \in K - \{0\}$,

$$\frac{dx_1(t)}{dt} = \lambda\, x_1\,(t) f_2(t, x_1(t), x_2\,(t-\tau_{12})) + x_1\,(t) m(\lambda - 1)$$

**(2.13)**

$$\frac{dx_2(t)}{dt} = \lambda\, x_2\,(t) f_2(t, x_1(t-\tau_{21}), x_2\,(t)) + x_2\,(t) m(\lambda - 1)$$

For all $t$ such that $x_1(t) \geq b_1^u/a_1^u$ *or* $x_2(t) \geq b_2^u/a_2^u$ we have $\frac{dx_1}{dt} < 0$ *or* $\frac{dx_2}{dt} < 0$; but these are impossible since $x_1, x_2$ are periodic in t. Thus by the continuity and periodicity of $x_1, x_2$ we have $x_1(t) \leq x_1^* < b_1^u/a_1^u$, $x_2(t) \leq x_2^* < b_2^u/a_2^u$ for all $t \geq 0$; by the definition and periodicity the same will follow for all $t \geq -\tau$. A similar argument holds when $x_1(t) < 0$ or $x_2(t) < 0$, noting that in such a case $dx_1/dt > 0$ or $dx_2/dt > 0$, again contradicting the periodicity of $(x_1, x_2)$ in *t*.

LEMMA 3. *The operator N maps* $K \cap (G_2/G_1)$ *in to K and is completely continuous*

PROOF. From the definition of N and (2.4) we derive that

$$N_1(t, x_1, x_2) \geq \int_0^\omega G_{11}\,(m, \omega, t-s) x_1(s)[m + b_{1l} - a_{1l} x_1(s)]ds$$

$$+ \int_0^\omega G_{12}\,(m, 0, t-s) x_1(s)[m + b_{1l} - a_{1l} x_1(s)]ds.$$

**(2.14)**

For $(x_1, x_2) \in K \cap (G_2/G_1)$, we have $0 < x_1 < b_1^u/a_1^u$ and hence we get $N_1(t, x_1, x_2) > 0$ for all $t \geq 0$ by a proper choice of the positive constant say $m_1$.

We make a similar choice say $m_2$ for *m* for the case $N_1(t, x_1, x_2) \geq 0$ and then choose *m* = max $(m_1, m_2)$. Thus we have

$$N: K \cap (G_2/G_1) \to K \qquad for\ t \geq 0,$$

And by the periodicity of N we have the same for $t \in -\tau$.

To prove the complete continuity of N we proceed as follows. We know that if $(x_1, x_2) \in B$ then $N(x_1, x_2) \in B$ and so we can restrict $t$ to the set $[0, n\,\omega], n\omega > r$. We write $N = L \circ Q$ where $L$ is the linear operator, $L : B \rightarrow B$, $L(u_1, u_2)(t)$

$$= \int_0^{\omega} G_{11}(m, \omega, t-s) \begin{bmatrix} u_1 & (s) \\ u_2 & (s) \end{bmatrix} ds, + \int_0^{t} G_{12}(m, 0, t-s) \begin{bmatrix} u_1 & (s) \\ u_2 & (s) \end{bmatrix} ds$$

**(2.15)**

And Q is the nonlinear operator

$$Q(u_1, u_2)(t) = \begin{matrix} g_1(m, t, u_1(t), u(t-\tau_{12})) \\ g_2(m, t, u_1(t-\tau_{21}), u_2(t)) \end{matrix} for\ \ t > 0,$$

**(2.16)**

$$Q(u_1, u_2)(t) = Q\,(u_1, u_2)(n\omega + t) \quad for\ t \in [-\tau, 0].$$

It can be found that

$$||L(u_1, u_2)(t)|| \leq max \left\{ \frac{m+b_1^u}{m}, \frac{m+b_2^u}{m} \right\} ||u_1, u_2||,$$

**(2.17)**

$$\left\| \frac{d}{dt} L(u_1, u_2) \right\| \leq max \left\{ \frac{m+b_1^u}{m}, \frac{m+b_2^u}{m} \right\} ||u_1, u_2||.$$

The complete continuity of the linear map $L$ follows from (2.17). Now since Q is continuous and maps bounded sets into bounded sets it will follow that $L \circ Q$ is completely continuous.

To prove the existence of periodic solution of (2.1)-(2.4) we will use the following result of Schmitt [9].

**THEOREM (Schmitt).** *Let $H_1$ and $H_2$ be nonempty bounded open neighborhoods of the origin in a Banach space B with a positive cone K such that $H_1 \subset H_2$ and let N: $K \cap (H_2 / H_1) \to K$ be completely continuous and which satisfies the following conditions:*

***(e)*** *There exists $(z_1, z_2) \in K \subset B$, $\|(z_1, z_2)\| = 1$ and a $\mu$ satisfying*

$$\mu > sup\{\|(x_1, x_2) - N(y_1, y_2)\|, (x_1, x_2) \in K \cap H_1, (y_1, y_2) \in K \cap \partial H_2\}$$ *such that all solution* $(x_1, x_2) \in K \cap (H_2/H_1)$ *of* $(x_1, x_2) = N(x_1, x_2) + \lambda(z_1, z_2) with\ 0 < \lambda < \mu\ satisfy\ (x_1, x_2) \notin \partial H_1.$

*(f) All solution* $(x_1, x_2) \in K \cap (H_2/H_1)\ of\ (x_1, x_2) = \lambda N(x_1, x_2), 0 < \lambda < 1\ satisfy\ (x_1, x_2) \notin \partial H_2.$

*Then N has at least one fixed point in* $K \cap (H_2/H_1)$.

We can now present our existence result.

**THEOREM 1.** *Assume that* $f_1, f_2$, *of* (2.1) *satisfy* (2.2)-(2.4). *then there exists at least one positive periodic solution of* (2.1)-(2.4).

PROOF. By Lemma 1, it is enough to show that the operator $N: K \cap (G_2/G_1) \to K \subset B$ has at least one fixed point in $G_2/G_1$ and we will do this by applying the result of Schmitt [9]. First we let

$$H_2 = \{(x_1, x_2) \in B | \ ||x_1, x_2|| < \max(b_1^u/a_1^u, b_2^u/a_2^u)\} \equiv G_2,$$

$$H_1 = \{(x_1, x_2) \in B | \ ||x_1, x_2|| < \min(b_{1l}/a_{1l}, b_{1l}/a_{1l})\} \equiv G_1.$$

By Lemma 3, $N: K \cap (H_2/H_1) \to K$ is completely continuous. Choose $\mu = \infty$ and $(z_1, z_2) \equiv (1.1)$. Then all solution $(x_1, x_2) \in K \cap (H_2/H_1)$ of $(x_1, x_2) = N(x_1, x_2) + \lambda(z_1, z_2), 0 < \lambda < \infty$, satisfy (2.1) and hence by their periodicity $(x_1, x_2) \notin \partial H_1$. The condition (a) of Schmitt's theorem is satisfied since $\mu = \infty$ by choice and $(x_1, x_2)$ and $(z_1, z_2)$ are periodic. By Lemma 2 the condition (b) of Schmitt's theorem is satisfied. Thus there exists at least

one fixed point of N in $K \cap (H_2/H_1)$ and hence there exists at least one positive periodic solution of (2.1)-(2.4) and the proof is complete.

We remark that so far no restriction on the magnitudes of the delays $\tau_{12}, \tau_{21}$ has been imposed except that they be nonnegative.

## 9. Uniqueness of the periodic solution

Applications of fixed point theorems to prove the existence of solutions of nonlinear problems usually suffer from the drawback that the uniqueness of solutions has to be proved often by other methods. We will use the following result of Smith [11] which is a slight variation of his Theorem 3.5 (see also Busenberg and Cooke [1]).

THEOREM (Smith). *Let N be an operator on a cone K in the Banach space B and let* $K^* = \{(x_1, x_2) \in K \cap B, 0 \leq x1,x2 \leq C*$ *where* $C*$*some positive constant is. Assume the following:*

(g) If $(x_1, x_2) \in K^* - \{(0,0)\}$ and $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

(h) N is monotonic nondecreasing on $K^*$.

(i) If $(x_1, x_2)$ is in the interior of $K^*$ and if $0 < \lambda < 1$ then there exists an $\eta > 0$ depending on $(x_1, x_2)$ and $\lambda$ such that $N(\lambda x_1, \lambda x_2) \geq \lambda(1+\eta) N(x_1, x_2)$. Then N has at most one fixed point different from $(0,0)$ in $K^* \subset B$.

The cone $K = K_1 \cap K_1$ where $K_1 = \{h \in B_\omega^\tau : h(t) \geq 0, t \in [-\tau, \infty)\}$ is slightly too wide in that solutions of the form $\{x_1(t), 0\}$ and $\{0, x_2(t)\}$ for (2.1) are also possible. We choose a narrower cone $K^*$ in $B$ as follows: define

$$S = \{(x_1, x_2) \in B : b_{1l}/a_{1l} \leq \|x_1\| \leq b_1^u/a_1^u,\ b_{2l}/a_{2l} \leq \|x_2\| \leq b_2^u/a_2^u,\}, \quad \textbf{(3.1)}$$

And let

$$K^* = \{\lambda(x_1, x_2) \in B : (x_1, x_2) \in S, \lambda \geq 0\}. \quad \textbf{(3.2)}$$

Since $\hat{G}_2/G_1 \subset K^*$, it is not difficult to see that N: $\hat{G}_2/G_1 \to B$ has at least one fixed point which is a periodic solution of (2.1)-(2.4). We will prove that this solution is unique by applying the result of Smith cited above. The following lemmas are needed first.

LEMMA 4. If $(x_1, x_2) \in K^* - \{0.0\}$ and if $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

PROFF. If $(x_1, x_2) = N(x_1, x_2)$ and $(x_1, x_2) \in K^* - \{0,0\}$, then $(x_1, x_2)$ is a periodic solution of (2.1) –(2.4) and hence we have

$$(b_{1l} - a_{1l}x_1)x_1 < \frac{dx_1}{dt} < (b_1^u - a_1^u x_1)x_1,$$
$$(b_{2l} - a_{2l}x_2)x_2 < \frac{dx_2}{dt} < (b_2^u - a_2^u x_2)x_2, \qquad (\mathbf{3.3})$$

And hence any nonzero periodic solution $(x_1, x_2)$ of (2.1)-(2.4) has to be such that

$$b_{1l}/a_{1l} < x_1(t) < b_1^u - a_1^u,$$

$$b_{2l}/a_{2l} < x_2(t) < b_2^u - a_2^u,$$

For $t > -\tau$, which proves the lemma.

LEMMA 5. The mapping $N$ defined by (2.5)-(2.9) is monotone nondecreasing on $K^*$.

PROOF. Let $(x_1, x_2) \in K^*$, $(y_1, y_2) \in K^*$ such that $(x_1 - y_1, x_2 - y_2) \in K^*$ and $(x_1, x_2) \neq (y_1, y_2)$. Then $x_1 \geq y_1$, $x_2 \geq y_2$. From the definition of $N$ we derive that

$N_1(t, x_1, x_2) - N_1(t, y_1, y_2)$

$$= \int_0^\omega G_{11}(m, \omega, t-s)\{m[x_1(s)] + x_1(s)f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s)f_1(s, y_1(s), y_2(s-\tau_{12}))\}ds$$

$$+ \int_0^t G_{12}(m, 0, t-s)$$

$$x\{m[x_1(S) - y_1(s)] + [x_1(s)f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s)f_1(s, y_1(s), y_2(s-\tau_{12}))]\}ds \qquad \mathbf{(3.4)}$$

With a similar relation for $N_2(t, x_1, x_2) - N_2(t, y_1, y_2)$. Since $f_1$ is continuous, $f$ is bounded on $K^*$; also it will follow from $(x_1, x_2) \neq (y_1, y_2)$ that

$$0 \not\equiv x_1(s) - y_1(s) \geq 0 \text{ and } 0 \not\equiv x_2(s) - y_2(s) \geq 0;$$

A sufficiently large positive constant $m$ will make the right side of (3.4) positive.A similar argument will lead to $N_2(t, x_1, x_2) - N_2(t, y_1, y_2) > 0$ perhaps with a larger constant $m$. Finally we can choose the bigger of these two constants $m$ to yield

$$N(x_1, x_2) - N(y_1, y_2) \text{ for } (x_1 - y_1, x_2 - y_2) \in K^* - \{(0,0)\}.$$

The monotonicity of $N$ on $K^*$ follows.

LEMMA 6. Let $(x_1, x_2) \in$ interior of $K^*$ and let $\lambda$ be $a$ constant such that $0 < \lambda < 1$. Suppose $f_1, f_2$ satisfy (2.2) - (2.4). Then there exists $a$ Positive constant $\eta$ such that

$$N(\lambda x_1, \lambda x_2) \geq \lambda(1+\eta) N(x_1, x_2). \quad \textbf{(3.5)}$$

PROOF. A direct computation leads to

$$N_1(t, \lambda x_1, \lambda x_2) - \lambda N_1(t, x_1, x_2)$$

$$= \lambda \int_0^{\omega} G_{11}(m, \omega, t-s) g_1\big(s, x_1(s), x_2(s-\tau_{12})\big) H_1(s, \lambda, x_1, x_2) ds$$

$$+ \lambda \int_0^{t} G_{11}(m, 0, t-s) g_1\big(s, x_1(s), x_2(s-\tau_{12})\big) H_1(s, \lambda, x_1, x_2) ds$$

**(3.6)**

Where

$$H_1(s, \lambda, x_1, x_2) = \frac{m\, x_1(s) + x_1(s) f_1\big(s, \lambda\, x_1(s), \lambda\, x_2(s-\tau_{12})\big)}{m\, x_1(s) + x_1(s) f_1\big(s, x_1(s), x_2(s-\tau_{12})\big)} - 1. \quad (\mathbf{3.7})$$

Since $f_1$ is decreasing (see (2.2)) with respect to $x_1$ and $x_2$ there exists a positive constant say $\eta_1$ such that

$$H_1(s, \lambda, x_1, x_2) \geq \eta_1 \text{ for } 0 < \lambda < 1, s > 0, (x_1, x_2) \in K^*, \quad \textbf{(3.8)}$$

And hence

$N_1(t, \lambda x_1, \lambda x_2) \geq \lambda(1+\eta_1)N_1(t, x_1, x_2)$.

A similar argument will show that there exists an $\eta_2 > 0$ such that

$N_2(t, \lambda x_1, \lambda x_2) \geq \lambda(1+\eta_2)N_2(t, x_1, x_2)$,

And the result follows if we choose $\eta = \min(\eta_1, \eta_2)$.

Now the following uniqueness result is a direct consequence of the result of Smith [11] and the Lemmas 4, 5, 6.

THEOREM 2. *Suppose $f_1, f_2$ of (2.1) satisfy (2.2) – (2.4). Then there exists a unique $\omega$ - periodic strictly positive periodic solution of (2.1).*

## 4. Linear stability of the periodic solution

Let $\{x_1^*(t), x_2^*(t), t \geq -\tau\}$ be the strictly positive unique periodic solution of (2.1). Let $\{x_1(t), x_2(t), t \geq -\tau\}$ be any other solution of (2.1). To study the linear stability of $(x_1^*, x_2^*)$ we let

$$\left.\begin{aligned} x_1(t) &\equiv x_1^*(t)\{1 + X_1(t)\}, \\ x_2(t) &\equiv x_2^*(t)\{1 + X_2(t)\}, \end{aligned}\right\} \quad t \geq -\tau, \tag{4.1}$$

And derive that the perturbations $(x_1, x_2)$ are governed by the linear variational system (after neglecting the non linear terms in $x_1, x_2$)

$$\left.\begin{aligned} \frac{dX_1(t)}{dt} &= -a_{11}(t)\, x_1^*(t)X_1(t) - a_{12}(t)x_2^*(t-\tau_{12})\, X_2(t-\tau_{12}) \\ \frac{dX_2(t)}{dt} &= -a_{21}(t)x_1^*(t-\tau_{21})x_1(t-\tau_{21}) - a_{22}(t)x_2^*(t)x_2(t), \end{aligned}\right\} \tag{4.2}$$

For $t > 0$, where

$$\left.\begin{aligned} -a_{11}(t) &= \frac{\partial f_1}{\partial x_1} \\ -a_{12}(t) &= \frac{\partial f_1}{\partial x_2} \end{aligned}\right\} \text{Evaluated at } (t, x_1^*(t), x_2^*(t-\tau_{12})), \tag{4.3}$$

$$\left.\begin{aligned} -a_{22}(t) &= \frac{\partial f_2}{\partial x_1} \\ -a_{22}(t) &= \frac{\partial f_2}{\partial x_2} \end{aligned}\right\}$$ Evaluated at $(t, x_1^*(t-\tau_{21}), x_2^*(t))$.

To establish the linear asymptotic stability of $(x_1^*(t), x_2^*(t))$ it is enough to prove that of the trivial solution of (4.2) - (4.3) We will do this by means of a combination of a comparison differential system and a Lyapunov function. First we define

$$\frac{d|X_1(t)|}{dt} = \begin{cases} \frac{dX_1(t)}{dt} & if\ X_i(t) \geq 0, \\ -\frac{d\setminus X_1(t)}{dt} & if\ X_i(t) < 0, \end{cases} \quad i = 1,2, \qquad \textbf{(4.4)}$$

And derive from (4.2)- (4.3) the following comparison system of differential inequalities:

$$\begin{aligned} \frac{d|X_1(t)|}{dt} &\leq -b_{11}, x_1^*(t)\,|X_1(t)| + b_{12}x_2^*(t-\tau_{12})\,|X_2(t-\tau_{12})|, \\ \frac{d|X_2(t)|}{dt} &\leq -b_{21}, x_1^*(t-\tau_{21})\,|X_1(t-\tau_{12})| - b_{22}x_2^*(t)\,|X_2(t)|, \end{aligned} \qquad \textbf{(4.5)}$$

For $> 0$, where

$$-b_{11} = \max_{t\in-\tau}\{-a_{11}(t)\},$$

$$\begin{aligned} b_{12} &= |\min_{t\in-\tau}\{-a_{12}(t)\}| \\ b_{21} &= |\min_{t\in-\tau}\{-a_{21}(t)\}| \end{aligned} \qquad \textbf{(4.6)}$$

$$-b_{22} = \max_{t\in-\tau}\{-a_{22}(t)\}$$

We can now formulate the result as follows. By the periodicity of $a_{ij}$, the constants $b_{ij}$ are independent of $\tau$ (see also (2.2)).

THEOREM 3. *The unique strictly positive periodic solution* $\{x_1^*(t), x_2^*(t), t \in -\tau\}$ of $(2.1) - (2.4)$ *is linearly asymptotically stable provided*

$$|-b_{11}| > b_{21} \text{ and } |-b_{22}| > b_{12},$$

**(4.7)**

*Where the constants* $b_{ij}$ $(\mathrm{i,j} = 1,2)$ *are as defined in* (4.6).

PROOF. Consider the scalar valued Lyapunov-like function $v(t) = V(t, X_{1,}X_{2,})$where

$$v(t) = V(t, X_{1,}X_{2,}) = |X_{1,}(t)| + |X_{2,}(t)| \int_{t-\tau_{12}}^{t} b_{12} x_2^*(u)|X_{2,}(u)|du$$

$$+\int_{t-\tau_{21}}^{t} b_{21} x_1^*(u)|X_{1,}(u)|\mathrm{du}. \qquad \textbf{(4.8)}$$

Since $x_1^*, x_2^*$ are strictly positive, it is found from (4.7)-(4.) that

$$\frac{dv}{dt} \leq -[(|b_{11}| - b_{21})\, x_1^*(t)|X_{1,}(t)| + (|b_{22}| - b_{22})\, x_2^*(t)|X_2(t)|] \qquad \textbf{(4.9)}$$

For $t > 0$. It follows from (4.7)-(4.9) that $V(t, X_{1,}X_{2,}) \geq |X_{1,}(t)| + X_{2,}(t)|$ and $dv/dt \leq 0$,which will imply the stability of the trivial solution of (4.2)-(4.3). We claim that (4.9) also implies the asymptotic stability of the trivial solution of (4.2)-(4.3)in the sense that

$$\lim_{t-\infty} [\,|X_{1,}(t)| + |X_{2,}(t)|\,] = 0,$$

**(4.10)**

For all bounded initial conditions $X_{1,}(s), X_{2,}(s), s \in \{\mathcal{T}, 0\}$. $Suppose (4.10) is\ not$ true Then we can find an increasing sequence $\{t_m\} \to \infty$ as $m \to \infty$ such that

$$(|b_{11}| - b_{21})\, x_1^*(t_m)|X_{1,}(t_m)| + (|b_{22}| - b_{12} x_2^*(t_m)|X_2(t_m)| > \varepsilon,$$

M=1, 2, 3…

**(4.11)**

Using the right derivative *dv/dt* of $v$ we derive that

$$v(t_m - \varepsilon) - \mathrm{v}(t_m) \leq -\int_{t_m-\varepsilon}^{t_m} \frac{\mathrm{dv\,(T)}}{\mathrm{dT}}\,\mathrm{d}\tau.$$

**(4.12)**

If $\varepsilon$ is small enough we have from (4.9)(4.11),(4.12)that $dv(\tau)/dr < (-\varepsilon/2)$ and

$$v(t_m - \varepsilon) - \mathrm{v}(t_m) \geq \varepsilon^2/2 \text{ and hence}$$

$$v(t_m) \geq v(t_m - \varepsilon) - (\varepsilon^2/2)$$

$$\geq v(t_m - 1)\text{-}(\varepsilon^2/2)$$

$$\geq v(t_m - 2)\text{-}(\varepsilon^2/2)$$

$$\geq v(t_0)\text{-m}\,(\varepsilon^2/2)$$

Thus for sufficiently large $m, v(t_m)$becomes negative which contradicts the Nonnegative nature of $v\ for\ t \geq 0$.

this contradiction shows that $|X_{1,}(\mathrm{t})| + |X_2(\mathrm{t})| \to 0$ as $\mathrm{t} \to \infty$ and the proof is complete.

As an example let us consider the case of two species Lotka- Volterra type competition in a periodic environment for for which we have

$$f_1\big(\mathrm{t}, x_1(\mathrm{t}), x_2(\mathrm{t} - T_{12})\big)1 \quad \left\{ \frac{x_1(\mathrm{t}) + a x_2(\mathrm{t} - T_{12})}{K_1(\mathrm{t})} \right\}$$

$$f_2(\mathrm{t}, x_1(\mathrm{t} - {}_{21}) x_2(\mathrm{t})) = 1 - \left\{ \frac{\beta x_1(\mathrm{t} - T_{21}) + x_2(\mathrm{t})}{K_2(\mathrm{t})} \right\}$$

Where $\alpha, \beta$ are positive constants and $k_{1,} k_{2,}$ are strictly positive continuous periodic functions with a common period Equation (4.6)and (4.7) for this special case lead to

$$\left|\operatorname{axm}_{t>T}\{-a_{11}(t)\}\right| = \max_{t>T}(1/K_1(t)) > \min_{t>T}\left(-a_{21}(\mathrm{t})\right)\Big| = \min_{t\in}, \beta/K_2(\mathrm{t})$$

$$\max_{t\in T}\{-a_{22}(t)\}\Big| = \max_{t\in T}(1/K_2(t)) > \min_{t\in T}\left(-a_{12}(\mathrm{t})\right)\Big| = \min_{t\in} \mathfrak{a}/K_1(\mathrm{t})$$

The inequalities reduce in the absence of temporal variations in $k_{1,}k_{2,}$ to the inequality 0< $\mathfrak{a}$,$\beta$<1, which is well known for the global asymptotic stability of the positive steady state corresponding to the above example in a constant systemic domain the above example shows also that delays in interspecific interactions have no influence on the local asymptotic stability of the unique periodic solution of the periodic ecosystem A number of cases of this type in a constant environment have been established recently by the author (Gopalsamy [6]).A numerical example is illustrated at the end.

**Consequences**

It can be found that all the conclusions on the existence uniqueness and stability of a periodic solution of (2.1)-(2.4) are valid when $OT_{12}$=0 or $T_{12}$=0 or both $T_{12}$=0 and $T_{21}$=0: thus it will follow that delays in interspecific competition do not destabilize an otherwise stable coexistence (oscillatory or nonoscillatory) of two competing types of interacting particles atoms molecules transactions interaction. In forthcoming work we will extend the analysis of the system (2.1)-(2.4) to a system of n-competing species with a hypothesis similar to that of (2.4) It is however an open problem to investigate the effects of delays in intraspecific competition in a periodic milieu

Pianka ([8],pages 139-140) remarks that “In real population rate of increase competitive abilities and carrying capacities do very from individual to individual with population density in space and time Indeed temporal variation in the environment may often allow coexistence by continually altering the competitive abilities of population inhabiting it Moreover time lags are doubt lessly of some importance in real population Finally a heterogeneous environment may allow real competitors to evolve divergent resource utilization patterns and to reduce interspecific competitive inhibition

To understand the mechanism of selection pressure against competitive inhibition we first write the Malthusian parameters $\mu_1$, $\mu_2$ of fitness for the two species as follows (see MacArthur[7],page 226-230):

$$\mu_1=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_1\left(t,x_1(t),x_2(t-T_{12})\right)$$

$$\mu_2=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_2(t,x_1(t-T_{21}),x_2(t))$$

Since $\left(x_1(t),x_2(t)\right)\rightarrow\left(x_1^*(t),x_2^*(t)\right)$ for large t we can write the parameters of finess in the from

$$\mu_1=f_1\left(t,x_1^*(t),x_2^*(t-T_{12})\right)$$

$$\mu_2=f_2\left(t,x_1^*(t-T_{12}),x_2^*(t)\right)$$

In these relations we can use

$$x_2^*(t-T_{12}),x_2^*(t)-T_{12}\frac{dx_2^*(t-T_{12}+\theta_1T_{12}}{dt}$$

$$x_1^*(t-T_{12}),x_1^*(t)-T_{21}\frac{dx_1^*(t-T_{21}+\theta_2T_{21}}{dt}$$

$\theta_1,\theta_2,\in(0,1)$ and rewrite $\mu_1,\mu_2$ in the form

$$\mu_1=f_1\left(t,x_1^*(T),x_2^*(t)-T_{12}\frac{dx_2^*(t-T_{12}+\theta_1T_{12}}{dt}\right.$$

$$\mu_2=f_2\left(t,x_1^*(T),-T_{21}\frac{dx_1^*(t-T_{21}+\theta_2T_{21}}{dt},x_2^*(t)\right)$$

$\theta_1,\theta_2,\in(0,1)$

It is seen that the presence of the delays $T_{12}$ and $T_{21}$ provides additional sources of variability in the finess parameters and thereby in the selection pressure against interspecific competition Since a necessary and sufficient condition for natural selection to act is differential (variable) rates of confederation growth we find that time delays coupled with periodicity of an domain may allow real to evolve divergent resource utilization patterns and to reduce

Interspecific competitive inhibition for some remarks relating to the significance of a periodic environment to sustain the coexistence of a large number of confederation types in the absence of delays we refer to a recent article (Gopalsamy)The following example has been numerically solved and the solutions are displayed graphically for two different values of the delay parameter Ƭ:

$\frac{dx_1(t)}{dt} x_1(t)[1 - x_1(t) + 0.4x_2(t - \mathcal{T})/k_1(t)]$

$\frac{dx_2(t)}{dt} x_2(t)[1 - \{0.8\, x_1(t - \mathcal{T}) + x_2(t)\}/k_2(t)]$

Where $k_1(t) = 20 + \sin\ (\pi t)$ and $k_2(t)$=2.5+sin (πt)

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There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies

systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational

content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification.

## GOPALSAMY'S MODEL APPIED TO COSMOLOGICAL PROBLEMS

## GLOBAL ASYMPTOTIC STABILITY IN AN ALMOST PERIODIC LOTKA- VOLTERRA SYSTEM

**Abstract**

**Gopalsamy studied** Sufficient conditions are obtained for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of a Lotka- Volterra system with almost periodic coefficients. Research was published Australian. Math Soc. B.27 (1986) 346-360. Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of

decay rate of perturbations. Sole intention, main objective and primary concern of the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are counted in the same vein as that of predator prey. Note that the model is applicable to each and every module. Quintessentially each parameter should have been given different name such as superscript 1,2, etc., as that has been done in the foregoing and a formal concatenation done at the end. Due to spatial restrictions we have not been able to do the line of thought and manner of analysis as has been done hereinbefore. This is the most important aspect of the model. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 5. Introduction

At the outset we assume that the system. This article is concerned with the derivation of a set of "easily verifiable" sufficient conditions for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of the Lotka- Volterra system

$$\frac{dx_1(t)}{dt} = x_i(t)\left\{b_1(t) - \sum_{j=1}^{n} a_{ij}\,(t)x_j(t)\right\};$$

$$t>t_0;\ t_0\ \in (-\infty,\infty);\ i = 1,2,\dots,n, \qquad (1.1)$$

where $b_1, a_{ij}\ (i,j = 1,2,\dots.,n)$ are nonnegative almost periodic functions defined for $t \in (-\infty,\infty)$.

In question satisfies Haimovici's equation based on age wise stratification given in the foregoing. Note that Haimovici's equations can be applied to Lotka-Volterra system term by term.

Let R and $R_n$ respectively denote the set of all real numbers and the $n$ - dimensional real Euclidean space; $R_n^+$ will denote the nonnegative cone of $R_n$. Since almost- periodic functions defined on R are bounded, we define the constants $b_i^l, b_l^u, a_{ij}^l, a_{ij}^u, (i,j = 1,2,\dots,n)$ by the following:

$$\inf_{t\in R} b_l(t) = b_i^l; \qquad \sup_{t\in R} b_l(t) = b_l^u, \qquad i,j = 1,2,\dots,n.$$

$$\inf_{t\in R} a_{lj}(t) = a_{lj}^l; \qquad \sup_{t\in R} a_{lj}(t) = a_{ij}^u.$$

We shall study the almost- periodic system (1.1) with the following assumption on the coefficients of (1.1);

$$b_l^l > 0; \qquad a_{ll}^l > 0, \tag{1.2}$$

$$b_l^l > \sum_{\substack{j=1 \\ j\neq i}}^{n} a_{lJ}^u (b_j^u / a_{iJ}^l). \tag{1.3}$$

Since the solutions of (1.1) corresponding to nonnegative initial conditions remain nonnegative subsequently, it will follow that

$$\frac{dx_1}{dt} \leq x_l\{b_l^u - a_{ll}^l x_i\}; \qquad t > t_0, i = 1,2,\dots,n, \tag{1.4}$$

As a consequence of which we will have

$$0 < x_l(t_0) \leq b_l^u / a_{ll}^l = x_i^u \Rightarrow x_i(t) \leq b_l^u / a_{Il}^l, \qquad i = 1,2,\dots,n; t \geq t_0. \tag{1.5}$$

(1.3) And (1.5) together lead to

$$\frac{dx_1}{dt} \leq x_l \left\{ b_l^u \sum_{\substack{j=1 \\ j\neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right) - a_{ll}^l x_i \right\}, \qquad i = 1,2,\dots,n; \tag{1.6}$$

Now (1.2), (1.3) and (1.6) lead to

$$x_i(t_0 \leq \left( b_i^l - \sum_{\substack{J=1 \\ J\neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right) \right) (1/a_{ll}^u) = x_i^l$$

$$\Rightarrow x_l(t) \to x_i^l \ for\ t \to t_0,\ i = 1, 2, \dots, n. \quad (1.7)$$

6. It is our aim to show that the mode together with other models stated in the foregoing together with haimovicis model is applicable to whole lot of consummative debit – credit systems in cosmology, astrophysics, finance, economics, sociology, anatomy computation. We will not however divide the system studied here into 9 components/ modules due to restrictions on space….

From the foregoing preparation we have the following:

LEMMA 1.1. *Assume that* $b_i, a_{ij} (i, j = 1, 2, \dots, n)$ *are scalar nonnegative almost periodic functions defined on* $\boldsymbol{R}$ *such that (1.2) and (1.3) hold. Then the set S defined by*

$$S = \{x = (x_1, x_2, \dots, x_n) \in \boldsymbol{R}_n | x_* \to x_l \to x^*; i = 1, 2, \dots, n\} \quad (1.8)$$

*Where*

$$x_* = \min_{1 \to n} x_l^l; \qquad x_* = \min_{1 \to n} x_i^u \quad (1.9)$$

*is invariant with respect to (1.1)*

**Appellation, cognomen, compellation, designation, denomination of the conducive, infelicitous and unfortutious** responses to environment change may be influenced by changes in available domain and milieu under consideration. An example is the following study: Authors describe the gravitational degrees of freedom of the Schwarzschild black hole by one free variable. Authors introduce an equation which we suggest to be the Schrödinger equation of the Schwarzschild black hole corresponding to this model. We solve the Schrödinger equation explicitly and obtain the mass spectrum of the black hole as such as it can be observed by an observer very far away and at rest relative to the black hole. Our equation implies that there is no singularity inside the Schwarzschild black hole, and that the black hole has a certain ground state in which its mass is non-zero. **Schrodinger equation of the Schwarzschild black hole Jarmo Mäkelä (Cambridge U E-Print: gr-qc/9602008 | PDF** Particles, axions, atoms, molecules, quarks, Majorana fermions, bosons et al in its confederation capacity and aggregation processes, type interactions and interactions between landscape dynamics. Following bears ample testimony, infallible observatory to the fact there might be changes in the quantum behaviour due to merger of black holes. Following example bears enough testimony, infallible observatory and impeccable demonstration and apodictic evidence as how some of the models could be applied to cosmological problems. This is attributable and ascribable to the fact that the classification criterion is based on the age.

**How gravity kills (e) Schrödinger's cat Theorists argue that warped space-time prevents (e) quantum superpositions of large-scale objects. Elizabeth Gibney 17 June 2015 Propositions are written in model able form and can easy be scheduled.**

(45) If the cat in Erwin Schrödinger's famous thought-experiment behaved according to quantum theory, it would be able to exist in multiple states at once: both dead and alive.

(46) Physicists' common explanation for why we don't (e) see such quantum superpositions — in cats or any other aspect of the everyday world — is interference from (e) the environment.

(47) As soon as a quantum object interacts with (e&eb) a stray particle or a passing field, it picks (e) just one state, collapsing (eb) into our classical, everyday view.

(48) But even if physicists could completely isolate (eb) a large object in (eb) a quantum superposition, according to researchers at the University of Vienna, it would still collapse into (e&eb) one state — on Earth's surface, at least.

(49) "Somewhere in interstellar space it could be that the cat has (e) a chance to preserve quantum coherence, but on Earth, or near any planet, there's (=) little hope of that," says Igor Pikovski. The reason, he asserts, is gravity.

(50) Pikovski and his colleagues' idea, laid out in a paper published in Nature Physics on 15 June1, is at present only a mathematical argument. But experimenters hope to test whether gravity really does collapse (eb) quantum superpositions, says Hendrik Ulbricht, an experimental physicist at the University of Southampton, UK. "This is a cool, new idea, and I'm up for trying to see it in experiments," he says. Assembling the technology to do so, however, may take as long as a decade, he says.

(51) Cinema-goers who saw the film Interstellar are already familiar with the basic principle behind the Vienna team's work. Einstein's theory of general relativity states that an extremely massive object causes clocks near it to (e) run more slowly because (e) its strong gravitational field stretches (e&eb) the fabric of space-time (which is why a character in the film aged only an hour near a black hole, while (e) seven years passed on Earth). On a subtler scale, a molecule placed nearer the Earth's surface experiences (e&eb) a slightly slower clock than one placed slightly further away.

(52) Because of gravity's effect on (e&eb) space-time, Pikovski's team realised that variance in a molecule's position will also influence (e&eb) its internal energy — the vibrations of particles within (eb) the molecule, which evolve (eb) over time.

(53) If a molecule were put in a quantum superposition of two places, the correlation between position and (e&eb) internal energy would soon cause (eb) the duality to 'decohere' to (e) the molecule taking just one path, they suggest.

(54) "In most situations decoherence is due to (e) something external; here it's as though the internal jiggling is interacting with (e&eb) the motion of the molecule itself," adds Pikovski.

**A practical limit**

(55) No one has yet seen this effect because other sources of decoherence — such as magnetic fields, thermal radiation and vibrations — are (=) typically much stronger and cause (eb) quantum systems to collapse long before gravity becomes an issue. But experimenters are keen to try.

(56) Markus Arndt, an experimental physicist also at the University of Vienna, already tests whether quantum superpositions can be observed for (e) large, though not cat-sized, objects (see Fattening up Schrödinger's cat).

(57) He sends large molecules through (e&eb) a matter-wave interferometer, a system that gives (eb) each molecule a choice of (e) two different paths.

(58) In the classical view, a molecule travels down just one path; a quantum molecule effectively passes through (e&eb) both routes at once and interferes with (e&eb) itself to create (eb) a characteristic wave-like pattern (see image, below).

| |
|---|
| (59) A quantum interference pattern built up from complex molecules known as phthalocyanines.(deleted due to spatial constraints) |
| (60) A similar set up could be used to (e) test for gravity's ability to destroy (e) quantum behaviour: by comparing a vertical interferometer, in which the superposition should soon decohere due to (e) the stretching of time in one path versus (e&eb) the other, with a horizontal set up, where (e) the superposition could remain. |
| (61) Arndt, who has tested the effect for (e) molecules of up to 810 atoms2, points out those large molecules, would be (=) good for testing the gravitational effect because (e) they contain (e) lots of particles contributing to (e) internal energy. |
| (62) **But not only would the researchers have to further suppress the external environment to reduce (e) other decoherence effects, they would also need to (e) increase the separation of the two paths from micrometres to metres, or else use molecules perhaps 1 million times more massive. "It's certainly very challenging," says Arndt**. |
| (63) If gravity's effect does limit (e) quantum behaviour on Earth, tests of quantum reality for (e) large objects may eventually have to (e) move into space, says Angelo Bassi, a physicist at the University of Trieste in Italy. |
| (64) "But from a deep, fundamental point of view, this is nothing new," he says. A gravitational field is merely another environment to interact with, so invoking it does not (e) explain whether quantum behaviour might lead to (eb) classical reality if gravity's influence were mitigated — for example, by doing the experiment in (eb) gravity-free space. |
| (65) Nor does the effect described by Pikovski and colleagues say anything about quantum gravity: a theory that would unite gravity and (e&eb) quantum mechanics in (eb) a single description, which many researchers are working towards. |
| (66) "It's an interesting effect, but it's still quantum physics applied to (e&eb) classical general relativity. In that way it doesn't change our picture of the world," adds Bassi. Nature doi:10.1038/nature.2015.17773 |

Integrated factors these influences because they deal with ambience, mise en scene change or confederation dynamics, but rarely both. Time series of (e) space time suitability models with spatially explicit stochastic models to explore factors that influence the viability of type viability and confederation feasibility under stable and changing scenarios. Methodology could be applied to interrogation propositions and problems which might subsume one with the other notwithstanding that there is a space of singularities and there is space for completion of present from the past and future. An example is Time image movement from Deleuze. We clearly face a stiff challenge in our attempt to proffer Deleuzian treatments of photography as our critical project brushes up against the fact that Deleuze does not invest photography with (e&eb)the creative and re-creative potential that he associates with (e&eb) both cinema and modern painting. In short, while he highly values other visual artistic forms, he seemingly presents photographic texts as stagnate documents or tools that produce certainty, organize bodies and desires, and iterate hackneyed ideas. He even uses photography as something of a foil to demonstrate the innovation of the cinema and the originality of modern painters. In Francis Bacon: The Logic of Sensation (1981), he notes how "photography has taken over the illustrative and documentary role, so that modern painting no longer needs to fulfill this function, which still burdened earlier painters" (10). He later explains that "photographs are ways of seeing, and as such, they are illustrative and narrative reproductions or representations ... . They are what are seen, until finally one sees nothing else." He treats the photo as an instrument for reproducing representations of reality—a device that iterates images until they are ossified as established stories, icons, or even stagnant perceptions. Indeed, like Bergson before him, Deleuze compares photography to what he identifies as a reductionist model of perception: "what we see, what we perceive, are photographs" (74).[4] The photo is, in effect, an always-already passé sensational experience that numbs our sensitivity to the ongoing vitality of life and provides us with an efficient

understanding of the dynamism of our world. But it is compelling that Deleuze theorizes both the quotidian nature and the liberatory potential of photography; he treats it as a common, everyday visual experience that, despite its aesthetic and technological limitations, emancipates the modern painter to explore new artistic opportunities. This special issue of Rhizomes repeatedly provides artistic examples and critical strategies that invite us to see photography as replete with the power to unburden and emancipate that Deleuze associates with both film and painting. Indeed, one way to think of this project of reading Deleuze against Deleuze might be to ask: if photography can free painting and painters to pursue novel creative opportunities, can it likewise free itself, photographic practitioners, and even viewers to create and re-create anew? **Rhizomes » Issue 23 (2012) » Michael Kramp Unburdening Life, or the Deleuzian Potential of Photography Michael Kramp Lehigh University** dmk209@lehigh.edu

**Existence of an almost- periodic solution**

DEFINITION 2.1. A function $f(t,x)$,where $f$ is an $m$- vector, $t$ is a real scalar and $x$ is an $n$- vector, is said to be almost periodic in $t$ uniformly with respect to $x \in X \subset R_n$, $if\ f(t,x)$ is continuous in $t \in R\ and\ x \in X$, and if for any $\varepsilon > 0$, it is possible to find a constant $l(\varepsilon) > 0$ such that in any interval of length $l(\varepsilon)$ there exists a $\tau$ such that the inequality

$$\left|\left|f\ (t+\ \tau,x) - \ f(t,x)\right|\right| = \sum_{l=1}^{m} |f_l\ (t+\ \tau,x) = \ f_l(t,x)|$$

$$< \varepsilon \qquad (2.1)$$

is satisfied for all $t \in (-\infty,\infty), x \in X$. The number $\tau$ is called an $\varepsilon -$ translation number of $f(t,x)$.

Let $\{\lambda_j\}$ denote the set of all real numbers such that

$$\lim_{T\to\infty} T^{-1} \int_0^T f(t,x)\exp(-i\,\lambda_j\,t)dt \not\equiv 0 \quad for\ x \in X; i = \sqrt{-1}. \qquad (2.2)$$

It is known that when the set X is separable, the set of numbers $\{\lambda_j\}$ in (2.2) is countable. The set $\{\sum_1^N n_j\ \lambda_j\}$ for all integers N and integers $n_j$ is called the module of $f(t,x)$.

The following result concerned with module containment will be used in the proof of our existence theorem below.

**MODULE CONTAINMENT THEOREM** (Yoshizawa [7], page 18). *Let $f(t,x)$ and $g(t,x)$ be almost periodic in t uniformly for $x \in D \subset \boldsymbol{R}_n$. If for any compact set $S \subset D$ and for any sequence of real numbers $\{\tau_k\}$ having a limit as $k \to \infty$ (the limit being finite or$\pm\infty$) for which $\{f(t+\tau_k,x)\}$ is uniformly convergent on $\boldsymbol{R} \times$*

$S, \{g(t+,x)\}$ *is also uniformly convergent on* $\boldsymbol{R} \times S$ *then the module of* $g(t+,x)$ *is contained in the module of* $f(t+,x)$.

We can now formulate our existence theorem.

**THEORM 2.1** *Assume that the almost-periodic parameters* $b_1, a_{1J}$ $(i,j = 1,2,\dots,n)$ *of (1.1) satisfy (1.2)-(1.3). Furthermore assume that there exists a positive number* $\mu$ *such that*

$$\inf_{1 \in \boldsymbol{R}} a_{11}(t) > \sum_{\substack{J=1 \\ J \neq 1}}^{n} (\sup_{t \in \boldsymbol{R}} a_{Jl}(t)) + \mu, \quad i = 1,2,\dots,n. \tag{2.3}$$

*Then the system (1.1) has a strictly positive (component wise) almost-periodic solution, say* $w(t) = w_1(t), \dots, w_n(t), t \in \boldsymbol{R}$, *whose module is contained in that of*

$$f(t,x) = \{f_1(t,x_1,\dots,x_n), f_2(t,x_1,\dots,x_n), \dots, f_n(t,x_1,\dots,x_n)\}$$

*Where*

$$f_l(t,x_1,\dots,x_n) = x_l \left\{ b_l(t) - \sum_{j=1}^{n} a_{lj}(t) x_J \right\},$$

$$i = 1,2,\dots n; (x_1,x_2,\dots,x_n) \in S. \tag{2.4}$$

PROOF. The set $S$ defined by (1.8)-(1.9) is compact in $\boldsymbol{R}_n$. Let $\{\tau_h\}$ be any sequence of real numbers such that $\tau_k \to \infty$ as $k \to \infty$ and

$$\sum_{l=1}^{n} |f_i(t+\tau_k, x_1, \dots, x_n) - f_l(t, x_1, \dots, x_n)| \to 0 \quad \text{as } k \to \infty \tag{2.5}$$

Uniformly in $(x_1, \dots, x_n) \in S$ for $t \in (-\infty, \infty)$. for any real number $\beta$let $k_0(\beta)$ be the smallest value of $k$ such that $\tau k_0 + \beta \geq t_0$. Since the invariance of $S$ implies that any solution of (1.1)

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), x_2(t+\tau_k, t_0, x_2(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\} \in S$$

For all $t \geq \beta$ and $k \geq k_0(\beta)$ whenever $\{x_1(t_0), \dots, n_n(t_0)\} \in S$, our first task in the proof is to show that the sequence

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Converges to a continuous bounded function say $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined for $t \in [\beta, \infty)$ such that $w_1(t) \in [x_*, x^*], t \in [\beta, \infty), i = 1, 2, \dots, n$ and the convergence is uniform on all compact subsets of $[\beta, \infty)$. Since $x_*$ and $x^*$ are independence of $\beta$, it is enough to show that

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Is a Cauchy sequence on compact subsets of $[\beta, \infty)$.

Let $U$ be any compact subset of $\beta, \infty)$., and let be any arbitrary positive number. Choose an integer $n_0(\varepsilon, \beta) \geq k_0(\beta)$ so large that for $m \geq k \geq n_0(\varepsilon, \beta)$

and all $t \in (-\infty, \infty)$

$$\sum_{i=1}^{n} |b_1(t+\tau_m) - b_1(t+\tau_k)| < (\mu\varepsilon x_*)/(4x^*), \tag{2.6}$$

$$x^* \sum_{i=1}^{n}\sum_{j=1}^{n} |a_{lj}(t+\tau_m) - a_{lj}(t+\tau_k)| < (\mu\varepsilon x_*)/(4x^*) \tag{2.7}$$

$$(2nx^*/x_*)\exp[-\mu x_*(\beta + \tau_k - t_0)] < (\varepsilon/2x^*)$$

(2.8)

Denote $x_i(t, t_0, x_1, (t_0))$ $by$ $x_1(t)$ brevity, and consider a function $v(s)$ defined by

$$v(s) = v\left(x(s), x(s+\tau_m+\tau_k)\right)$$

$$= \sum_{i=1}^{n} |\log x_1, (s) - \log x_1(s+\tau_m+\tau_k)|, \qquad x \geq t_0, \tag{2.9}$$

Where $x(t) = \{x_1(t), \dots, x_n(t)\}$ denotes a solution of (1.1) with $\{x_1, (t_0), \dots, x_n(t)\} \in S$. the invariance of S for (1.1), together with the elementary mean value theorem of differential calculus, implies

$$(1/x^*)\sum_{l=1}^{n}|x_l(s)-x_1(s+\tau_m+\tau_k)|\ \upsilon(s)$$

$$(1/x^*)\sum_{i=1}^{n}|x_l(s)-x_1(s+\tau_m+\tau_k)|. \tag{2.10}$$

Calculating the right derivative $D^+\upsilon$ of $\upsilon$ along the solutions of (1.1) and simplifying

$$D^+\upsilon(s)=\sum_{i=1}^{n}\{sign[x_i(s)-x_i(s+\tau_m-\tau_k)]\}$$

$$\left\{\frac{d[logx_i(s)]}{ds}-\frac{d[logx_i(s+\tau_m-\tau_k)]}{ds}\right\}$$

$$\sum_{l=1}^{n}|b_l(s)-b_l(s+\tau_m-\tau_k)|$$

$$+\sum_{l=1}^{n}(-a_{il}(s)|x_i(s)-x_i(s+\tau_m-\tau_k)|+\sum_{\substack{J=1\\ j\neq l}}^{n}a_{ij}(s)|x_j(s)-x_j(s+\tau_m-\tau_k)|)$$

$$+\sum_{i=1}^{n}\sum_{j=1}^{n}|a_{ij}(s)|a_{ij}(s+\tau_m-\tau_k)|\,x_j(s+\tau_m-\tau_k) \tag{2.11}$$

$$-\mu x_*\upsilon\big(x(s),x(s+\tau_m-\tau_k)\big)+\sum_{l=1}^{n}|b_l(s)-b_i(s+\tau_m-\tau_k)|$$

$$+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|. \tag{2.12}$$

Let us set $s-\tau_k=\sigma$; then we have from (2.6)-(2.7) that

$$\sum_{i=1}^{n}|b_l(s)-b_1(s+\tau_m-\tau_k)|+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|.$$

$$= \sum_{l=1}^{n} |b_l(\sigma + \tau_k) - b_i(\sigma + \tau_m)|$$

$$+x^* \sum_{i=1}^{n} \sum_{J=1}^{n} |a_{ij}(\sigma + \tau_k) - a_{ij}(\sigma + \tau_m)|$$

$$< (\mu\varepsilon x^*)/(2x^*). \quad (2.13)$$

Now an integration of (2.12) over $(t_0, t - \tau_k)$ together with (2.13) leads to

$$v\big(x(t+\tau_k), x(t+\tau_m)\big) \leq v\big(x(t_0), x(t_0 + \tau_m - \tau_k)\big) \exp[-\mu x_*(t + t_k - t_0)]$$

$$+[(\mu\varepsilon x_*)/(2x^*)]/\mu x_*)$$

$$\leq [(2x_n^*)/x_*]\exp[-\mu x_*(t + t_k - t_0)] + (\varepsilon/2x^*)$$

$$< (\varepsilon/x^*) \text{ On using (2.8), (2.10)} \quad (2.14)$$

For $m \geq n_0(\varepsilon, \beta)$ and for all $t \in \cup \subset [\beta, \infty]$, Thus we have from the definition of $v$ and (2.14) that

$$\sum_{l=1}^{n} |log x_i(t + \tau_k) - log x_l(t + \tau_m)| < (\varepsilon/x^*) \quad (2.15)$$

And hence using $x_* \leq x_i(t) \leq x^*, i = 1, 2, \dots, n; t \geq t_0$,

$$\sum_{i=1}^{n} |x_i(t + \tau_k) - x_i(t + \tau_m)| < \varepsilon \text{ for all } t \in \cup \subset [\beta, \infty], m \geq n_0(\varepsilon, \beta). \quad (2.16)$$

The existence of a function $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined on $[\beta, \infty]$, such that $w_i(t) \in [x_*, x^*], i = 1, 2, \dots, n; t \in [\beta, \infty]$ follows from (2.16) and the fact that $x_i(t) \in [x_*, x^*], i = 1, 2, \dots, n;$ for $t \geq t_0$. Since $\beta$ is arbitrary, $w$ is in fact

Defined on $(-\infty, \infty)$ and hence we have

$x_i(t+\tau_k) \to w_l(t)$ as $k \to \infty$ uniformly on compact subsets of $(-\infty,\infty)$, $i = 1,2,\dots,n.$

(2.17)

Our next task is to show that $w: R \to S$ is differentiable and is a solution of (1.1). Since $x(t) = \{x_1(t),\dots,x_n(t)\}$ satisfies (1.1) we have $\sum_{i=1}^{n} |\frac{d}{dt} x_i(t+\tau_k) - \frac{d}{dt} x_i(t+\tau_m)|$

$$\sum_{i=1}^{n} |f_l\left(t+\tau_k, x_1(t+\tau_k),\dots, x_n\left(t+\tau_k\right)\right)$$

$$- f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)| \qquad (2.18)$$

$$+\sum_{l=1}^{n} |f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$- f_l\left(t+\tau_m, x_1(t+\tau_m),\dots,x_n(t+\tau_m)\right)|$$

Where $f_l\ (i = 1,2,\dots,n)$are as in (2.4). Since $x(t+\tau_k) = x_1(t+\tau_k),\dots, x_n(t+\tau_k) \in S$ for large $\tau_k$ for $t$ in each compact subset of $(-\infty,\infty)$ there exists an integer $n_1(\varepsilon) > 0$ such that if $m, k \geq n_1(\varepsilon)$, then for any $\varepsilon > 0$,

$$\sum_{i=1}^{n} |f_l\left(t+\tau_k, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$-f_1\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)| < (\varepsilon/2). \qquad (2.19)$$

Similarly since $x_l(t+\tau_k) \in S$ and $x_l(t+\tau_m) \in S$ for $i = 1,2,\dots,n$ there is an integer $n_2(\varepsilon) > 0$ such that for large enough $m, k$, such that $m, k \geq n_2(\varepsilon)$, we have

$$\sum_{l=1}^{n} |f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$-f_1\left(t+\tau_m, x(t+\tau_m),\dots,x_n(t+\tau_m)\right)| < (\varepsilon/2) \qquad (2.20)$$

This equation will be very much useful in chemistry if used with the modules mentioned and written in concatenation with each other. It should be borne in mind that the applicability is modulewise and in no case concatenation should be marginalized. In all cases parameters are different which could be represented with different superscripts and subscripts like 1,23,…….etc., **Anal Bioanal Chem. 2004 Sep; 380(1):36-45. Epub 2004 Jul 29. Electronic tongue for pharmaceutical analytics: quantification of tastes and masking effects Legin A et al**

The organoleptic aspects of pharmaceutical formulations affect their acceptability to the patient and hence can have an important effect on concordance with treatment. Objective evaluation of these aspects, particularly the taste of the formulation and the drug substance it contains, is difficult. Whilst volunteer taste panels can be used to good effect their utility is limited, particularly during very early stage development when the toxicological profile of the active pharmaceutical ingredient (API) is yet to be established in detail. A potentiometric "electronic tongue" has been applied to analyse a variety of 41 individual substances and mixtures of particular interest for pharmaceutical research and development. The electronic tongue (ET) was capable of discriminating between substances with different taste modalities and could also distinguish different substances eliciting the same basic taste; the ET is promising in terms of quantifying the content of each substance and has an ability to detect nuances of the basic taste (e.g. lingering or short-lived). After calibration the electronic tongue was successfully applied to predicting bitterness strength of binary mixtures with a sweetener in terms of "apparent" or "perceived" quinine content. In order to render a formulation palatable it is often necessary to mask the (usually bitter) taste of the API by the addition of masking agents such as sweeteners and flavours. The ET proved capable of distinguishing between formulations with different levels of sweetener and/or flavour in a manner that was consistent with their masking efficiency as perceived by a small human taste panel. A suitably calibrated ET could have the benefit of providing the pharmaceutical formulator with reliable data concerning the taste of the product quickly and with a reduced need to ask volunteers to taste active pharmaceutical samples. Early development activities could be facilitated when human tasting is usually not possible in the absence of the required toxicological data.

Thus if $m \neq n_0(\varepsilon) = max\{n_1, n_2\}$, we have

$$\sum_{l=1}^{n} |\frac{d}{dt}\, x_1(t+\ \tau_k) - \frac{d}{dt} x_1\ (\ (t+\tau_m)| < \varepsilon \qquad (2.21)$$

Which shows that

$$\lim_{k\to\infty} \frac{d}{dt} x_1(t+\ \tau_k) \quad exists\ for\ i = 1, 2, \dots, n$$

Uniformly on all compact subsets of $(-\infty, \infty)$. Thus we have

$$\lim_{k\to\infty} \lim_{k\to\infty} \frac{x_i(t+\tau_k+h) - \ x_i(t+\tau_k)}{h}$$

$$= \lim_{h\to 0} \lim_{k\to 0} \frac{x_l(t+\tau_k+h) - \ x_l(t+\tau_k)}{h} \qquad (2.22)$$

$$= \lim_{h\to 0} \frac{w_1(t+h) - \ w_1(t)}{h}, \qquad i = 1, 2, \dots, n,$$

Showing that (since the left side of (2.22) exists) $w_i (i = 1, 2, \dots, n)$ is differentiable. Furthermore we have from (2.22) that

$$\frac{dw_l(t)}{dt} = \lim_{k\to\infty} [\, f_i\ \big(t+\tau_k, x_1(t+\tau_k), \dots, x_n\ (t+\tau_k)\big)$$

$$-f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big)$$

$$+ - f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big)]$$

$$= \ -f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big); \qquad i = 1, 2, \dots, n, \qquad (2.23)$$

By the uniform almost periodicity of $f_1, f_2, \dots, f_n$ on $R \times S$ and hence $w(t) = \big(w_1(t), \dots, w_n(t)\big)$is a solution of (1.1).

In order to complete the proof, we have to show that $w(t) = \big(w_1(t), \dots, w_n(t)\big)$ is an almost- periodic function whose module is contained in that of $\{f_1(t, x_1, \dots, x_n), f_2(t, x_1, \dots, x_n), \dots, f_n(t, x_1, \dots, x_n)\}$ for $(x_1, x_2 \dots, x_n) \in S$. For this purpose it is in fact sufficient (by the module containment theorem) to show that for any sequence of real numbers $\{\tau_k\}$ for which

$$\{f_1(t + \tau_k, x_1, \dots, x_n), f_2(t, + \tau_k, x_1, \dots, x_n), \dots, f_n(t, + \tau_k, x_1, \dots, x_n)\}$$

Converges uniformly on R for all $(x_1, \dots, x_n) \in S$, the sequence $\{w_1(t, + \tau_k) \dots, w_n(t + \tau_k)\}$ converges uniformly on R where $\tau_k$ tends to a finite number or $\pm\infty$ as $k \to \infty$.

In the case where $\tau_k \to \tau \neq \pm\infty$ as $k \to \infty$, we have $w_l(t, + \tau_k) \to w_i(t + \tau), i = 1, 2, \dots, n$ uniformly on R as $k \to \infty$. Let us suppose $\tau_k \to \infty$ as $k \to \infty$. For any $\varepsilon > 0$ there exists an integer $n_0(\varepsilon)$ such that if $m$ [illegible] $n_0(\varepsilon)$ then

$$\sum_{l=1}^{n} |b_l(t + \tau_k) - b_l(t + \tau_m)| < (\varepsilon/4x^*)x_*, \qquad (2.24)$$

$$x^* \sum_{i=1}^{n} \sum_{j=1}^{n} |a_{ij}(t + \tau_k) - a_{ij}(t + \tau_m)| < \left(\frac{\varepsilon}{4x^*}\right)\mu x_* \qquad (2.25)$$

$$n(2x^*/x_*) \exp[\mu x_* \tau_k] < (\varepsilon/2x^*). \qquad (2.26)$$

Consider the function $v\big(w(s), w(s + \tau_m - \tau_k)\big)$ for $s \in [t, t + \tau_k]$ defined by

$$v\big(w(s), w(s + \tau_m - \tau_k)\big) = \sum_{l=1}^{n} |\log w_l(s) - \log w_l(s + \tau_m - \tau_k)|. \qquad (2.27)$$

Using the fact $w: \boldsymbol{R} \to S$ and $w$ is a solution of (1.1), calculate the right derivative $D^+v \ of \ v$;

$$D^+v\big(w(s), w(s + \tau_m - \tau_k)\big) = \sum_{i=1}^{n} \{\text{sign}[w_l(s) - w_i(s + \tau_m - \tau_k)]\}.$$

$$\left\{\frac{dw_i(s)}{ds} \cdot \frac{1}{w_l(s)} - \frac{dw_l(s + \tau_m - \tau_k)}{ds} \frac{1}{w_l(s + \tau_m - \tau_k)}\right\}$$

$$\leq -\mu \sum_{l=1}^{n} |w_l(s) - w_l(s+\tau_m-\tau_k)| + \sum_{l=1}^{n} |b_i(s) - b_l(s+\tau_m-\tau_k)|$$

$$+x^* \sum_{l=1}^{n}\sum_{J=1}^{n} |a_{lJ}(s) - a_{lJ}(s+\tau_m-\tau_k)| \quad (2.28)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+\sum_{i=1}^{n} |b_l(s) - b_l(s+\tau_m-\tau_k)| \quad (2.29)$$

$$+x^* \sum_{l=1}^{n}\sum_{j=1}^{n} |a_{lJ}(s) - a_{ij}(s+\tau_m-\tau_k)|.$$

If we set $s-\tau_k = \sigma$ then we have from (2.29),

$$D^+ v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+\sum_{l=1}^{n} |b_l(\sigma+\tau_k) - b_l(s+\tau_m)|$$

$$+x^* \sum_{l=1}^{n}\sum_{j=1}^{n} |a_{ij}(\sigma+\tau_k) - a_{ij}(\sigma+\tau_m)| \quad (2.30)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+[(\varepsilon\mu x_*)/(4x^*)] + [(\varepsilon\mu x_*)/(4x^*)]$$

(2.31)

Integrating (2.31) over $(t, t+T_k)$

$$v(w(t+T_k)w(t+T_m))$$

Another MODEL where the model is directly applicable with the scheduled ones are the case of insurance and corporate finance. Insurance: Mathematics and Economics Volume 42, Issue 1, February 2008, Retrieval of Black–Scholes and generalized Erlang models by perturbed observations at a fixed time Dedicated to Professor André Dubey on the occasion of his 60th birthday Daniel Neuenschwander

S-stable laws on the real line (more generally on Hilbert spaces), associated with some non-linear transformations (so-called "shrinking operations"), were introduced in [Jurek, Z.J., 1977. Limit distributions for truncated random variables. In: Proc. 2nd Vilnius Conference on Probability and Statistics, June 28–July 3, 1977. In: Abstracts of Communications, vol. 3, pp. 95–96; Jurek, Z.J., 1979. Properties of s-stable distribution functions. Bull. Acad. Polon. Sci. Sér. Math. XXVII (1), 135–141; Jurek, Z.J., 1981. Limit distributions for sums of shrunken random variables. Dissertationes Math. vol. CLXXXV]. In [Jurek, Z.J., Neuenschwander, D., 1999. S-stable laws in insurance and finance and generalization to nilpotent Lie groups. J. Theoret. Probab. 12 (4), 1089–1107], the authors interpreted s-stable motions on the real line as limits of total amount of claims processes (up to a deterministic premium) of a portfolio of excess-of-loss reinsurance contracts and showed that they led to Erlang's model or to Brownian motion. In [Neuenschwander, D., 2000b. On option pricing in models driven by iterated integrals of Brownian motion. In: Mitt. SAV 2000, Heft 1, pp. 35–39], we considered stochastic integrals whose integrand and integrator are both independent Brownian motions, thus modelling a stochastic volatility; as a result we got an analogue of the Black–Scholes formula in this model, confirming a result of Hull and White [Hull, J., White, A., 1987. The pricing of options on assets with stochastic volatility. J. Finance XLII (2), 281–300]. In the present paper, we will look at a common generalization of these processes, namely s-stable motions on the real line perturbed by a stochastic integral whose integrand and integrator are both (not necessarily independent) s-stable motions. The main result will be that if we can observe the distribution of such so-perturbed s-stable motions (together with the values of the perturbing processes) at time t=1, then we can identify the whole model (including the perturbation) among all models with Lévy processes perturbed by an iterated stochastic integral of two Lévy processes (in the Gaussian case) resp. among all models with a compound Poisson process with drift perturbed by an iterated stochastic integral of two compound Poisson processes (in the completely non-Gaussian case if the perturbing processes have no drift) without knowing anything about the history or about its distribution during $0 \leq t < 1$. This applies, e.g., to a situation

where several assets obey the same model and one can estimate the distribution at time one by looking at the values of all these assets at time t=1. Interestingly enough, it will be convenient to treat the whole matter in the algebraic framework of the so-called Heisenberg group. This is a concept coming in fact from quantum mechanics and is in a certain sense the simplest non-commutative Lie group.

$\leq e^{-\mu x^* T_k} v\big(w(t), w(t+T_m-T_k)\big) + (\varepsilon/2x^*)$

$\leq e^{-\mu x^* T_k} \sum_{i=1}^{n} |\log\ w_i(t) - \log\ w_i(t+T_m-T_k)| + (1/x_*) + (\varepsilon/2x^*)$

$\leq n(2x^*/x_*)\exp[-\mu x_* T_k] + (\varepsilon/2x^*)$

$<(\varepsilon/x^*)$.

Since

$$v\big(w(t+T_k)w(t+T_m)\big) \leq \sum_{i=1}^{n} |\ w_1(t+T_k) - w_1(t+T_m))/x^*)$$

We have from(2.33)that

$$\sum_{i=1}^{n} |\ w_1(t+T_k) - w_1(t+T_m)| < \varepsilon\ for\ m \geq k \geq n_0(\varepsilon)$$

Which show that the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\}$converges uniformly on **R**. By a similar argument the above conclusion is also valid when $T_k \to -\infty\ as\ k \to \infty$. Thus the uniform convergence of the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\}\ as\ k \to \infty\ on$ **R** follows and this completes the proof of the theorem.

**COROLLARY.** Suppose $b_i a_{ij}$ (i, j = 1.2, … , n)in (1.1)are periodic with a common period say $T > 0$. *Then if* (1.2), (1.3)*and* (2.3)*hold the system* (1.1)*will have a periodic solution* $w(t) = \{w,(t), \ldots\ldots, w_n(t)\}$*of period* $T$ *such that*

$$w_1(t+T) = w_1(t), \qquad t \in \boldsymbol{R}: i = 1,2, \ldots, n.$$

**PROOF.** Proof follows from that of Theorem 2.1 if we Choose $T_k = KT(K = 0,1,2, \ldots)$*and note that*

$$x_1(t+KT) - w_1(t) \to 0\ as\ k \to \infty; i = 1,2, \ldots, n; t \in \boldsymbol{R},$$

$$x_i(t+(K+1)t) - w_1(t) \to 0 \ \ as \ \boldsymbol{K} \to \infty; i = 1,2,\dots,n; t \in \boldsymbol{R},$$

$$x_i(t+T+KT) - w_1(t+T) \to 0 \ \ as \ k \to \infty; i = 1,2,\dots,n; t \in \boldsymbol{R},$$

Implying that

$$w_1(t+T) = w_i(t); \quad i = 1,2,\dots,n; t \in \boldsymbol{R},$$

## 3. Global asymptotic stability

By definition we say that an almost-periodic solution say $w(t) = \{w_1(t),\dots,w_n(t)\} of (1.1) is\ globally\ asymptotically$ stable (or attractive) if and only if every other solution x(t)={ $x_1(t),\dots,x_n(t)$}with $x_1(o) > 0, i = 1.2,\dots,n$ is defined for all t□0 and satisfies

$$\lim_{t\to\infty} |w_1(t) - x_1(t)| = 0; \quad i = 1,2,\dots,n.$$

A consequence of such a global asymptotic stability is that there cannot be another strictly positive almost-periodic solution of (1.1)

**THEOREM 3.1** Suppose the almost- Periodic coefficients $b_1$, $a_{ij}$ (i, j = 1,2, … , n)

Satisfy (1.2)and (1.3)and the following

$$\inf_{t\in R} a_{ij}(t) > \left\{ \sum_{\substack{j=1 \\ j\neq 1}}^{n} \sup a_{ij}(t) \right\} + \mu; \qquad i = 1,2,\dots,n$$

For some positive number $\mu\ Then$ (1.1)has an almost-periodic solution (with strictly positive components) which is globally asymptotically stable.

We will omit the proof since the proof is identical with that of Theorem 3.1 of [3]

Under condition (3.2) the almost –periodic solution $w(t) = \{w_1(t), \ldots . w_n(t)\}$ of (1.1)enjoys a somewhat strong from of stability even if the coefficients of (1.1)are subjected to a class of perturbations We note (see Yoshizawa [7]page17)that if a function b:R→ $R$ is almost periodic then H(b),the "hull of b" is defined as follows:

H(b)= g|g:**R**→ $\boldsymbol{R}$ $and\ g(t)\lim_{k\to\infty} b(t+T_k) for$ some sequence

$T_k$ of real numbers the limit being uniform in $t \in R$.

By Theorem 2.2 of Yoshizawa [7, page 10], members of the hull of almost periodic functions are themselves almost periodic.

A system of the from

$$\frac{dy_1(t)}{dt} = y_1(t)\left\{B_i(t) - \sum_{J-1}^{n} A_{ij}\ (t)y_1(t)\right\},$$

*i=1,2,…,n; t>*$t_0 \in R$,

Where

$$B_i \in H(b_i); (A_{ij}), \qquad \text{i, j} = 1,2, \ldots, \text{n},$$

is known as a system from the hull of (1.1) As a consequence of the definition of the hull in (3.3)and (3.5)we shall have

$$\inf_{t\in R} B_i(t) = \text{b}_i^{\,l}; \quad \sup_{t\in R} B_i(t) = \text{b}_i^{\,l};$$

$$\inf_{t\in R} A_{ij}(t) = \text{b}_i^{\,l}; \quad \sup_{t\in R} A_{ij}(t) = \text{a}_{ij}^{\,u};$$

and hence the coefficients of(3.4)- (3.5)satisfy the same conditions as (1.2),(1.3)and (3.2)It will then follow from theorems(2.1)and (3.1)that solutions of (3.4)-(3.5)with $y_1(t_0) > 0, i = 1,2, \ldots, n\ will\ have$ the following property:

$$y_1(t) > 0 \ \ for\ t \geq t_0; i = 1,2,\dots,n,$$

$$\sum_{i=1}^{n} y_1(t) remains\ bounded\ for\ t \geq t_0$$

**DEFINITION** A Solution $x(t) = \{x_1(t), \dots, x_n(t)\}, t \in t_0 \in \boldsymbol{R}$,of (1.1)is said to be stable under perturbations from the hull of (1.1)if the following holds :let $y(t) = \{(y_1(t), \dots, y_n(t)\}$ be any solution of (3.4)-(3.5)such that $y_1(t_0) > 0$, i=1,2,…,n,and for some positive constants $\beta_1, \beta_2 \beta_1, \leq y_1(t) \leq \beta_2, t \geq t_0, \ \ i = 1,2,\dots,n$ Then given any $\varepsilon > 0$,there exists a constant $\delta > 0$ such that

$$\sum_{j=1}^{n} |x_1(t_0) - y_1(t_0)| < \delta$$

$$\sum_{i=1}^{n} \sup_{t \in R} |b_1(1) - B_1(t)| < \delta \quad = \sum_{i=1}^{n} |x_1(t) - y_i(t)| < \varepsilon$$

$$\sum_{i=1}^{n} \sum_{j=1}^{n} | \quad \sup_{t \in R} |a_{ij}\ (t) - A_{ij}(t)| < \delta$$

For $t \geq t_0$.

We can now formulate the following result which shows that all solution of (1.1)are stable under perturbations from the hull of (1.1)

**THEOREM 3.2**. Assume that the conditions of Theorem (3.1) hold. Then every strictly positive solution of (1.1) ( including its unique almost-periodic solution) is stable under perturbations from the hull of (1.1).

PROOF. Let $x(t) = \{x_1(t), \dots, x_n(t)\}$ and $y(t) = \{y_1(t), \dots, y_n(t)\}$ be any two strictly positive solutions of (1.1) and (3.4) – (3.5) respectively for $t \geq t_0 \in R$. Let $\beta_1, \beta_2$ be constants such that

$\beta_1 \leq x_i(t) \leq \beta_2$ ,
$\beta_1 \leq y_i(t) \leq \beta_2$ , $i = 1,2,\dots,n; t \geq t_0$ (3.7)

Consider now a function $V(t) = V(x(t)y(t))$ defined by

$$V(x(t)y(t)) = \sum_{i=1}^{n} |\log x_1(t) - \log, y_i(t)|, \quad t \geq t_0 \in \text{R}. \tag{3.8}$$

It is easy to see that

$$\frac{1}{\beta_2} \sum_{i=1}^{n} |x_i(t) - y_i(t)| \leq V(x(t)y(t))$$

$$\leq \frac{1}{\beta_1} \sum_{i=1}^{n} |x_i(t) - y_i(t)|, \; t \geq t_0 \in \text{R}. \tag{3.9}$$

Calculating the right derivative $D^+V$ of $V$ and using (1.1) and (3.4) we derive (after simplification) that

$$D^+V\left(x(t), y(t)\right) \leq -\mu \sum_{i=1}^{n} |x_i(t) - y_i(t)| + \sum_{i=1}^{n} |b_i(t) - B_i(t)|$$

$$+\beta_2 \sum_{j=1}^{n} \sum_{i=1}^{n} |a_{ij}(t) - A_{ij}(t)|; \quad t \geq t_0,$$

$$\leq -\mu\, \beta_1\, V\left(x(t), y(t)\right) + \sum_{i=1}^{n} |b_i(t) - B_i(t)|$$

$$+ \beta_2 \sum_{i=1}^{n} \sum_{j=1}^{n} |a_{ij}(t) - A_{ij}(t)|; \quad t \geq t_0, \tag{3.10}$$

An integration of (3.9) over $[t_0, t]$ with an application of (3.8) leads to

$$\frac{1}{\beta_2}\sum_{i=1}^{n}|x_i(t)-y_i(t)| \leq V(x(t),y(t)) \leq V(x(t_0),y(t_0))$$

$$+\left[\sum_{i=1}^{n}\sup_{t\in R}|b_i(t)-B_i(t)|+\beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}|a_{ij}(t)-A_{ij}(t)|\right]\frac{1}{\mu\beta_1},$$

$$t \geq t_0,$$

And hence

$$\sum_{i=1}^{n}|x_1(t)-y_1(t)| \leq \left(\frac{\beta_2}{\beta_1}\right)\sum_{i=1}^{n}|x_1(t_0)-y_1(t_0)|$$

$$+\frac{\beta_2}{\mu\beta_2}\left[\sum_{i=1}^{n}\sup_{t\in R}|b_i(t)-B_i(t)|\right.$$

The construction reason is because is criticized in a number of usage guides: The reason for the long delays was because the costs greatly exceeded the original estimates. One objection to this construction is based on its redundancy: the word because (literally, by cause) contains within it the meaning of reason; thus saying the reason is because is like saying "The cause is by cause," which would never be said. A second objection is based on the claim that because can introduce only adverbial clauses and that reason is requires completion by a noun clause. Critics would substitute that for because in the offending construction: The reason for the long delays in completing the project was that the costs.

$$\left.+\beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}|a_{ij}(t)-A_{ij}(t)|\right]. \qquad (3.11)$$

Now for any $\varepsilon > 0$ if one chooses a $\delta > 0$ such that

$$\delta < \min\left\{\frac{\varepsilon}{3}\frac{\beta_1}{\beta_2}, \frac{\varepsilon}{3}\frac{\beta_1}{\beta^2}\mu\right\} \qquad (3.12)$$

Then (3..6) will be satisfied and this completes the proof

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# . Jia Li's Prescientiial Dynamics And Contential Staticity Of Accentuation Attrition Models Applied To Cosmological And Astrophysical And Quantum Mechanical Engineering Problems

**Abstract**: **Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics ET AL ). We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems. Predator prey model literally form bastion, pillar, post, stylobate, sentinel towards the end of dismasting many of the asymptotic and solutional behaviour in finance, economic, and as we shall see in cosmology, quantum computation and other fields.**

## MODEL FORMULATION

Let $p_1(t, a_1)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_1}{dt} = -b_1 Q_1 + g_1(g_1)Q_1$$

Where $b$ is the natural death rate in the absence of credit types and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_1(0) = 0, \quad g_1(P_1) > 0, \; P_1 \; 0, \qquad \lim_{P_1 \to \infty} g_1(P_1) = (K_1)_1 \; \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_1 + (\rho_a)_1 + D_1(a_1, P_1, Q_1)\rho_1 + 0, \qquad 0 < a_1 \; \infty,$$

$$\rho_1(t, 0) = \int_0^{\infty} n_1(a_1, P_1, Q_1)\rho_1(t_1, a_1)da_1$$

Where $D_1$ is the mortality function and $n_1$ is the fertility function.

Put $_1$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_1}{dt}\right]_1 = x_1\left[\frac{dQ_2}{dt}\right]_1 - x_2\left[\frac{dQ_1}{dt}\right]_1$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_1 - x_4\left[\frac{dQ_2}{dt}\right]_1$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_1 - x_6\left[\frac{dQ_3}{dt}\right]_1$$

Where $x_1, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_1$ and $g_1(P_1)Q_1$

Now when you add the term $(b_1Q_1)$ with $g_1(P_1)(Q_1)$ we are infect see is that $(b_1Q_1)$ is credited incrementally part of $g_1(P_1)(Q_1)$ or instant nearly, when subtracted from $\frac{dQ_1}{dt}$, it means that $g_1(P_1)(Q_1)$ is slowly eaten by $\frac{dQ_1}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_1 = \mu_1(a_1) + (f_1)_1(P_1)_1 + (f_2)_1(Q)_1 \quad (2)$$

$$n_1 = \beta_1(a_1)h_1(Q_1), \quad (3)$$

Where $\mu_1$the age- dependent debit rate whit is $\mu_1(a_1)$ 0 $\forall a > 0$; $(f_1)_1$ is the density- dependent death rate with $f_1(0) = 0$, $(f_1^1)_1(P_1)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_1)_1(P_1) = (k_1)_1$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_1(0) = 0, (f_2^1)_1(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_1(Q_1) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^1(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_1 + (\rho_a)_1 = -\mu_1(a_1) + (f_1)_1P_1) + \big(f_2)_1(Q_1)\big)(\rho_1)_1 \quad (4)$$

$$\rho_1(t_1, 0) = h_1(Q_1)\int_0^\infty \beta_1(a_1)\rho_1(t, a_1)da, \quad (5)$$

$$\frac{dQ_1}{dt}$$

$$= -b_1Q_1 + g_1(P_1)Q_1 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_1$ and $Q_1$ of (4-6) satisfy

$$0 \leq \rho_1(t, a_1) \leq \rho_1^*(t, a_1),$$

$$0 \leq Q_1(t) \leq Q_1^*(t),$$

Where $\rho_1^*(t, a_1)$ and $Q_1^*(t)$, are solutions of

$$(\rho_t^*)_1 + (\rho_a^*)_1 = -(\mu_1(a_1) + (f_1)_1(P_1^*) + (f_2)_1 Q_1^*))\, \rho_1^*, \quad (7)$$

$$\rho_1^*(t, 0) = \int_0^\infty \beta_1(a_1)\, \rho_1^*(t, a_1) da_1 \quad (8)$$

$$\frac{dQ_1^*}{dt} = -b_1 Q_1^* + g_1(P_1^*) Q_1^*, \quad (9)$$

With $P_1^*(t) = \int_0^\infty \rho_1^*(t, a_1) da_1$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_1^*(t)}{dt} = P_1^*(t)\big(r_1 - (f_1)(P_1^*) - (f_2)_1(Q^*)\big), \quad (10)$$

$$\frac{dQ_1^*(t)}{dt} = Q_1^*(t)(-b_1 + g_1(P_1^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_1(a) e^{-\int_0^a \mu_1(\tau) dt} e^{-ra}\, da_1 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_1^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(t)dt}\, da_1 > 1, \qquad (H1)$$

Which guarantees $r_1 > 0$.

*Remark.* If $\beta_1(a) = [\beta_0]_1 a_1 e^{-\alpha_1 a_1}, \mu_1(a_1) \equiv \mu_1$ then the solution of (12) is

$$r_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_1]_1(P_1) = 0, [f_2](Q_1) = \lambda_1 Q_1\ h_1(Q_1) = 1, \qquad and\ \ g_1(P_1) = c_1 P_1$$

$$[\rho_1]_1 + [\rho_a]_1 = -(\mu_1 + \lambda_1 Q_1)\rho_1$$

$$\rho_1(t,0) = \int_0^\infty [\beta_0]_1 a_1 e^{-\alpha_1 a_1}\rho_1(t,a_1)\, da_1$$

$$\frac{dQ_1}{dt} = -b_1 Q_1 + c_1 P_1 Q_1$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_1}{dt} = \omega_1 P_1 - \lambda_1 P_1 Q_1$$

$$\frac{dQ_1}{dt} = b_1 Q_1 - c_1 P_1 Q_1$$

Where $\omega_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$ THE EQUILIBRIUM $(\rho_0)_1(a_1), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_1(a_1), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_1 + (\rho_a)_1 = -\big(\mu_1(a_1) + (f_1)_1(P_1)\big)\rho,$$

$$\rho_1(t,0) = \int_0^\infty \beta_1(a_1)\rho_1(t,a_1)da_1 \qquad (13)$$

The solution of (13) has the form of

$$\rho_1(t,a_1) = Ce_1^{r(t-a_1)} - \int_0^a (\tau)d\tau - \int_0^1 (f_1)_1\big(P_1(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_1(t) = C_1 e^{r_1 t - \int_0^t (f_1)_1(P_1(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_1^{-1}(r)$ if $k_1 > r$.

This is equivalent to

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{(k_1)_1 a_1} da_1 < 1.$$

In fact, if we define

$$R_1(s_1) = \int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau)d\tau} e^{-s_1 a_1}\, da_1$$

$R(s)$ is a decreasing function. $R_1(r_1) = 1$ And $R_1(k_1)_1 < 1$ imply $r_1 < (k_1)_1$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_1(a_1), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{k_1 a_1} da_1 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_1 \equiv \int_0^\infty (\rho_0)_{11}(a_1)da_1$ *satisfies*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{-(f_1)_1(P_0)a_1} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_1(a_1), 0)$

Let $u_1(t, a_1) = \rho(t, a_1) - \rho_0(a_1)$ Then

$$N_1(t) \equiv \int_0^\infty u_1(t, a_1)da_1 = \int_0^\infty \rho_1(t, a_1)da_1 - \int_0^\infty (\rho_0)_1(a_1)da = P_1(t) - P_0)_1$$

And

$$(u_1)_1 + (u_a)_1 = (\rho_t)_1 + (\rho_a)_1 - (\rho_0^1)_1(a_1)$$

$$= -\mu_1(a_1) + (f_1)_1(P_1) + (f_2)_1(Q_1))(u_1 + (\rho_0)_1) + \mu_1(a_1)(\rho_0)_1 + (f_1)_1(P_0)_1(\rho_0)_1$$

$$= -\mu_1(a_1)u_1 - (f_1)_1(P_1) + (f_2)_1(Q_1))u_1 - (\rho_0)_1(f_1)_1(P_1) - (f_1)_1(P_0)_1$$

$$+ (f_2)_1(Q_1)), \qquad (15)$$

Where $(\rho_0^1)_1(a_1)$ denotes $d(\rho_0)_1(a_1)/da_1$

Expansion $(f_1)_1(P_1)$ and $(f_2)_1(Q_1)$ at $(P_0)_1$ and 0 yield respectively

$$(f_1)_1(P_1) = (f_1)_1((P_0)_1 + N_1)$$

$$= (f_1)_1(P_0)_1 + (f_1^1)_1(P_0)_1N_1 + o_1(N_1), \qquad (16)$$

$$(f_2)_1(Q_1)$$

$$= (f_2^1)_1(0)Q + o_1(Q_1) \qquad (17)$$

Write $f_1(P_0) = f_{10}, f_1'(P_0) = f_{10}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_1)_1 + (u_a)_1 \approx -(\mu_1(a_1) + (f_{10})_1)u - (f_{10}^1)_1(\rho_0)_1(N_1) - (f_{20}^1)_1(\rho_0)_1(Q_1)$$

$$u_1(t,0) = \rho_1(t,0)(\rho_0)_1(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t,a) + (\rho_0)_1(a_1)da_1 - (\rho_0)_1(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t,a_1)da_1 + (h_1(Q_1) - 1)(\rho_0)_1(0) \qquad (18)$$

Expanding

$$h_1(Q_1) = 1 + (h^1)_1(0)Q_1 + 0(Q_1) \equiv 1 + h_0'Q_1 + 0Q_1$$

And substituting it into (18), we arrive at

$$u_1(t,0) \approx \int_0^\infty \beta_1(a_1)u_1(t,a_1)da_1 + (h_{a_1}')(\rho_0)_1Q_1$$

The linearization is the following:

$$[u_1] + [u_1] = -(\mu_1(a_1) + [t_{10}]_1u_1 - [\alpha_1]_1(a_1)N_1 - [\alpha_2]_1(a_1)Q_1$$

$$u_1(t,0) = \int_0^\infty \beta_1(a_1)u_1(t,a_1)da_1 + [\xi_1]Q_1 \qquad (19)$$

$$\frac{dQ_1}{dt} = -b_1Q_1$$

where

$$[\alpha_1]_1(a_1) = [t_1^1 a]_1[p, o]_1(a_1)d2(a_1) = [f_{20}^1]_1(\rho\, o\,)_1(a_1)\xi = [\mathrm{h\ o}^1]_1[\rho \mathrm{o}]_1(\mathrm{o}) \text{ and } \mathrm{b}_1 = \mathrm{b}_1 - g_1(P_1)$$

If $g_1(P_0) > b$, clearly, $(\rho_0)_1(a_1), 0)$ is unstable and the predator population is persistent. Assume $g_1(P_0)_1 < b$.

Then, $Q_1(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_1 + u_a = -(\mu_1(a_1) + (f_{10})_1)u_1 - (\alpha_1)_1(a_1)N_1$$

$$= -(\mu_1(a_1) + f_{10})u_1 - \alpha_1(a_1)\int_0^\infty u_1(t, a_1)da_1$$

$$u_1(t, 0) = \int_0^\infty \beta_1\,(a_1)u_1(t, a_1)da_1$$

Using the method developed by Cushing [2, 3], we suppose

$$u_1(t, a_1) = \tilde{\mathrm{u}}_1(a_1)^{e^{\gamma(t-a_1)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_1}\, da_1$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_1)$ gives,

$$\tilde{\mathrm{u}}_1(a_1) = e^{-\int_0^a \mu_1(\tau_1)d\tau - (f_{10})_1 a}\left(\tilde{\mathrm{u}}_1(0) - w\int_0^a e^{\int_0^\tau \mu_1(s_1)ds + (f_{10})}\tau_1\,(\alpha_1)_1(\tau_1)e^{\gamma_1\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_1(a_1)e^{-\gamma_1 a_1 - (f_{10})_1 a_1 - \int_0^a \mu_1(\tau_1)d\tau}\right)\tilde{\mathrm{u}}_1(0)$$

$$+w\int_0^{\infty}\beta_1(a_1)e^{-\gamma_1 a_1-(f_{10})_1 a_1-\int_0^a \mu_1(\tau_1)d\tau}\int_0^a(\alpha_1)_1(\tau_1)e^{-\gamma_1\tau_1}e^{(f_{10})_1\tau_1+\int_0^{\tau}\mu_1(s_1)ds}d\tau_1\,da_1$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty}e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}da\right)u(0)$$

$$+\left(-1-\int_0^{\infty}e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}\int_0^a e^{f_{10}\tau+\int_0^{\tau}\mu(s)ds+y\tau}\alpha_1(\tau)d\tau\,da\right)w=0.\ (23)$$

Define $Y(\gamma,a)=e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta\,(a)Y(\gamma,a)da\right)u(0)$$

$$+w\int_0^{\infty}\beta(a)Y\,(\gamma,a)\int_0^a\alpha_1(\tau)Y^1(\gamma,\tau)d\tau\,da=0,(24)$$

$$\left(\int_0^{\infty}Y(\gamma,a)da\right)u(0)$$

$$+\left(-1-\int_0^{\infty}Y(\gamma,a)\int_0^a\alpha_1(\tau)Y^{-1}(\gamma,\tau)d\tau\,da\right)w=0.$$

Equation (20) has a solution of the form of $u(t,a)=u(a)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $u(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I(y,a)=\int_0^a\rho_0(\tau)Y^{-1}\,(\gamma,\tau)d\tau$. We obtain the characteristic equation

$$f_{10}^{l}\left(\int_0^{\infty} Y(\gamma,a)da\right)\left(\int_0^{\infty}\beta(a)Y(\gamma,a)I(\gamma,a)da\right)$$

$$+\left(1-\int_0^{\infty}\beta(a)Y(\gamma,a)da\right)\left(1+f_{10}^{l}\int_0^{\infty}Y(\gamma,a)I(\gamma,a)da\right)=0. \quad (25)$$

Since

$$\rho_0(a)=\rho_0(0)e^{-\int_0^a \mu(\tau)d\tau - f_{10}a}$$

$$I(\gamma,a)=\int_0^a \rho_0(\tau)e^{\gamma\tau+f_{10}\tau+\int_0^{\tau}\mu(s)ds}d\tau$$

$$=\frac{\rho_0(0)}{\gamma}(e^{\gamma a}-1)$$

Let

$$E(a)=e^{-f_{10}a-\int_0^a \mu(\tau)d\tau}$$

Then $[Y_1]_1[\gamma_1]_1,a_1)=\ E_1(a_1)e^{-\gamma_1 a_1}$. by a simple calculation,

$$\int_0^{\infty}\beta_1(a_1)Y_1(\gamma_1,a_1)I(\gamma_1,a_1)da=\frac{[\rho_0]_1(0)}{\gamma}\left(1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}\ da\right);$$

$$\int_0^{\infty}Y_1(\gamma_1,a_1)I(\gamma_1,a_1)da=\frac{[\rho_0]_1(0)}{\gamma_1}\left(\int_0^{\infty}E_1(a_1)e^{-\gamma_1 a_1}\ da_1\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_1(0)}{\gamma_1}[f_{10}^{1}]\int_0^{\infty}E_1(a_1)e^{-\gamma_1 a_1}da_1\ (\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1)$$

$$+\left(1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right).$$

$$\left(1+\frac{(\rho_0)_1(0)}{\gamma_1}(f_{10}^{'})_1\left(\int_0^{\infty}E_1(a_1)da_1-\int_0^{\infty}E_1(a_1)e^{-\gamma_1 a_1}da_1\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right)(1+\frac{(f_{10}^{'})_1}{\gamma_1}(\rho_0)_1(0)\int_0^{\infty}E_1(a_1)da_1)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_1(a_1)Y_1(\gamma_1 a_1)da_1=1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\neq 0,$$

Since $\int_0^{\infty}\beta_1(a_1)Y_1(\gamma_1 a_1)\int_0^{a}\alpha_1(\tau_1)Y^{-1}(\gamma_1,\tau_1)d\tau_1 da_1>0$. Therefore, the characteristic solution is

$$\gamma_1=-(f_{10}^{'})_1(\rho_0)_1\int_0^{\infty}E_1(a_1)da_1.$$

Since $f_1'(P_0)_1>0$, we obtain the following result

THEOREM 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_1),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_1),0)$ which is related to the extinction of debit summation as has been referred to in the foregoing and to The possible persistence of the prey

population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_1 + (\mu_1(a_1) + [f_1]_1(P_1^*) + [f_2]_1(Q_1^*))P_1^* = 0, \quad (26)$$

$$p_1^*(0) = h_1(Q_1^*)\int_0^\infty \beta_1(a_1)p_1^*(a_J)p_1^*(a_1)da_1 \quad (27)$$

$$g_1(P_1^*) - b_1 = 0, \quad (28)$$

Where $P_1^* = \int_0^\infty p_1^*(a_1)\, da_1$

Equation (28) yields $P^* = g^{-1}(b)$,

From (26),

$$p_1^*(a) = p_1^*(0)e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}$$

Substituting into (27), we need to solve

$$\left(1 - h_1(Q_1^*)\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}da_1\right)p_1^*(0) = 0$$

For $Q_1^*$

Fix $p_1^*(0)$ and $P_1^*$ Let

$$H_1(z_1) = 1 - h_1(z_1)\int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1}da_{,1}$$

And assume that

$$[k_3]_1 \int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(g_1^{-1}(b_1))a_1 - [k_2]_1 a_1}\, da_1 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_1(z_1)] > 0.$$

Clearly,

$$\frac{dH_1}{dz_1} = h_1(z)(\,[f_2^{'}]_1(z_1)\,) \int_0^{\infty} a_1\beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - f_1(P_1^*)a_1 - f_2(z_1)a_1} da_1$$

$-h^{'}(z_1)\int_0^{\infty} \beta_1(a_1)\, e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1}\ da_1$

By the assumptions on $\left[f_2^1\right]_1 and\ [h^1]^1 dH_1/dz > 0$. Hence, if

$$H_1(0) = 1 - \int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1} da_1 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_1(0) < 0$ follows if $g_1^{-}(b_1) < [f_1^{-1}]_1(r_1)$, where $r_1$ is the solution of (12).

Indeed, since

$$R_1(s_1) = \int_0^{\infty} \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 e^{-y_1^a}} da_1$$

Is a decreasing function of $s_1$ when $P_1^* = [g^{-1}]_1(b_1) < [f_1^{-1}]_1(r_1)$, that is $[f_1]_1(P_1^*) < r$, $R\big(f_1(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_1(g^{-1}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1\left(g_1^{-1}(b)\right)a_1} > 1. \qquad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_1^*(a_1), Q_1^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_1^*(a_1), Q_1^*)$ is more complicated. We again consider perturbations $m_1(t_1 a_1) = p_1(t_1, a_1) - p_1^*(a_1)$ and $W_1(t) = Q_1(t) - Q^*$.Then

$$M_1(t) \equiv \int_0^\infty m_1(t, a_1)da_1 = P_1(t) - P_1^*,$$

$$[m_1]_1 + [m_a]_1 = [p_1]_1 + [p_a]_1 - \left[p^*\right]_1(a_1) \approx -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1$$

$$-[f_{1*}^1]_1 p_1^*(a_1)M - [f_{2*}^1]_1 p_1^*(\mathrm{a})W_1,$$

$$m_1(t,0)p_1(t,0) - p_1^*(0) \approx h_* \int_0^\infty \beta_1(a_1)m_1(t,a_1)\,da + \frac{[h_*^1]_1}{[h_*]_1}\, p_1^*(0)W_1$$

$$\frac{dW_1}{dt} = \frac{dQ_1}{dt} \approx [g_*^1]_1 Q_1^* M_{,1}$$

Where $[f_{1*}]_1$ denoting $[f_1(P^*)]_1$, $[f_{1*}]_1$ denoting $[f_2(P^*)]_1$, $\left[f_{1*}^{'}\right]_1$ denoting $\left[f^{'}{}_1(P^*)\right]_1 \left[f_{2*}^{'}\right]_1$ denoting

$[f_2]_1\,(Q_1^*,\quad h_{1*}^{'}$ denoting $h_*^{'}$ denoting $h_1^{'}\,(Q_1^*)$, and $\left[g_*^{'}\right]_1$

Denoting $\left[g^{'}\right]_1(P_1^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_1]_1 + [m_a]_1 = -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1 - [f'_{1*}]_1[p_1^*](a_1)M_1 - [f'_{2*}]_1 p_1^*(a_1)W_1$$

$$m_1(t,0) = [h_*]_1 \int_0^\infty \beta_1(a_1)m_1(t,a_1)da_1 + \frac{[h_*^1]_1}{[h^*]_1}[p^*]_1(0)\,W_{,1} \qquad (30)$$

$$\frac{dW_1}{dt} = [g_*^1]_1[Q_1^*]M_1$$

Integrating $[m_1]_1(t,a_1)$ with respect to a from 0 to $\infty$, we have

$$\frac{dM_1(t)}{dt} = \int_0^\infty [mt]_1(t,a_1)da_1$$

$$= -\int_0^\infty [m_a]_1\,(t,a_1)da_1 - \int_0^\infty \mu(a_1)m(t,a_1)da_1$$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f'_{1*}]_1 \int_0^\infty p_1^*(a_1)da_1\right)M_1 - [f_{2*}^1]_1 \int_0^\infty p_1^*(a_1)da_1 W_1$$

$$= m_1(t,0) - \int_0^\infty \mu_1(a_1)\,m_1(t,a_1)da_1$$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f_{1*}']_1 P_1^*\right)M_1 - [f_{2*}']_1 P_1^* W_{,1}$$

From (30), we arrive at the equations

$$\frac{dW_1(t)}{dt} = \int_0^\infty ([h_*]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1$$

$$-\left([f_{1*}] + [f_{2*}]_1 + [f_{1*}']_1 P_1^*\right)M_1 + \left(\frac{[h_*']_1}{[h_*]^1}\,P_1^*(0) - f_{2*}' P_1^*\right)W_1,$$

(31)

$$\frac{dW_1(t)}{dt}\left[g^{'}{}_{*}\right]_1 Q_1^{*} M_{.1}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{*}]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1 = ([h_{*}]_1\beta_1(\theta_1) - \mu_1(\theta_1))M_{,1}$$

Where $0< \theta_1 = \theta_1(t) < \infty$.

And if $a_1^2\,[a_1^2]_1<4\,[a_2]_1<4\,a_2 e$,

Re $\lambda_1\left[\lambda_1\right]_1 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_1 = -\frac{1}{2}a_1[a2\ ]_1 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_{*}\,[h_{*}]_1\beta_1(a_1) - \mu_1(a_1)\right|$. Then$|n_1(t)|$ ⊠Δ∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium($\rho^{*}(a), \boldsymbol{Q}^{*}$) is asymptotically stable if either $a_1^2 <4\,a_2 e$ and $[a_1] > 2\Delta$,or $[a_1^2]$ ⊠4$[a_2]$e and $(a_1 - \sqrt{(a_1^2)_1 - 4(a_2)_1 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+1}\right)$ and $\beta_1(a_1) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_1(a_1)\mathrm{e}^{-\int_0^{\mathrm{a}} \mu_1(\tau_1)+d\tau_1 \mathrm{da}} = \frac{1}{(\mathrm{m}+\mu_0+1)^2} + \frac{2}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_1 + [\mu_0]$ 0.5. Assume$[f_1](x_1)= K_1 x_1$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_1(x_1)= [0x_1]/(1 + x_1)$; the functional responses $[[f_2](x) = [0]\, x_1/(1 + x_1)$and $h_1(x_1)=(1+\omega_1 x_1)/(1 + x_1)$(See Holling [7] and Rashevsky [13])

Equation (14),now has the form of

$$\frac{m_1 + [\mu_0] + [\mathrm{K}][P_0])}{[m_1]_1 + (\mu_0)_1 + [\mathrm{K}][P_0])_1^3\ [\mathrm{K}]_1 [P_0]_1^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_1$i e $[P_0] < b/(\delta_1 - b_1)$then the equilibrium$(P_0)\ [P_0]_1(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_1$

$[f_{1*}] = K_1\ \mathrm{P}^*\ [f_{2*}] = \frac{\sigma \mathrm{Q}_1^*}{1+\mathrm{Q}_1^*} \quad [h_*] = \frac{1+\omega_1 \mathrm{Q}_1^*}{1+\mathrm{Q}_1^*}$

$[f_{1*}] = K_1\ [\mathrm{f}_{2*}^1] = \frac{\sigma_1}{(1+\mathrm{Q}_1^*)^2} [\ \mathrm{g}_*^1] = \frac{\sigma_1}{(1+\mathrm{P}_1^*)^2} \quad [\mathrm{h}_*^1] = \frac{\omega_1 - 1}{(1+\mathrm{Q}_1^*)^2}$

## AN ANALOGY BETWEEN PREDATOR PREY AND ACCENTUATION ATTRION MODELS

**Best example how the model could be applied is the Schrödinger's equation itself wherein RHS gobbles up LHS, and the terms on RHS do the same with each other. It is like adding milk to water. Whether it happens instantaneously or time lag is there is a different question altogether dealt at the opportune points**

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| **Best example that can be given is the formulation of the quantum wave equation albeit in a nascent form for the recently found gravitational waves. This has been studied in length in various papers of gravitational waves, wherein the attribution of how the model is used is described. Note the calculation holds for only one module. This is in consideration to the fact that the relation is beteen the both debit systems and credit systems which as we have enucleated in the foregoing bears ample testimony for the application of the model** |
| **Gravitational quantum wave equation:** |
| **Equation Time-dependent equation (Wikipedia)** |
| **(1) In quantum mechanics, the Schrödinger equation is (=) a partial differential equation that describes how the quantum state of a quantum system changes with (e&eb) time. It was formulated in late 1925, and published in 1926, by the Austrian physicist Erwin Schrödinger.[1]** |
| **(2) In classical mechanics Newton's second law, (F = ma), is used to make (eb) a mathematical prediction as to what path a given system will take following a set of known initial conditions.** |
| **(3) In quantum mechanics, the analogue of Newton's law is (=) Schrödinger's equation for (e) a quantum system (usually atoms, molecules, and subatomic particles whether free, bound, or localised).** |
| **(4) It is not a simple algebraic equation, but in general a linear partial differential equation, describing (eb) the time-evolution of the system's wave function (also called a "state function").[2]:1–2** |
| **(5) The concept of a wavefunction is (=) a fundamental postulate of quantum mechanics.** |
| **(6) Although Schrödinger's equation is often presented as a separate postulate, some authors [3]: Chapter 3 show that some properties resulting from Schrödinger's equation may be deduced just from (e) symmetry principles alone, for example the commutation relations.** |
| **(7) Generally, "derivations" of the Schrödinger equation demonstrate (eb) its mathematical plausibility for describing (eb) wave–particle duality, but to date there are no universally accepted derivations of (e) Schrödinger's equation from (e) appropriate axioms.** |
| **(8) In the Copenhagen interpretation of quantum mechanics, the wave function is (=) the most complete description that can be given of (e) a physical system.** |
| **(9) Solutions to Schrödinger's equation describe (eb) not only molecular, atomic, and subatomic systems, but also macroscopic systems, possibly even the whole universe.[4]:292ff** |

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| **(10) The Schrödinger equation is consistent with (e&eb) classical mechanics but not with (e&eb) special relativity.** |
| **(11) Making quantum mechanics consistent with (e&eb) special relativity is one of the great achievements of (e) Quantum Field Theory and requires (e) the introduction of the creation and annihilation operators so that (e) the number of particles can no longer be considered constant.** |
| **(12) The Schrödinger equation is not the only way to make predictions in (eb) quantum mechanics—other formulations can be used, such as Werner Heisenberg's matrix mechanics, and Richard Feynman's path integral formulation** |
| **(13) The form of the Schrödinger equation depends on (=) the physical situation (see below for special cases).** |
| **(14) The most general form is the time-dependent Schrödinger equation, which gives a description of a system evolving with time: [5]:143** |
| **(15) A wave function that satisfies the non-relativistic Schrödinger equation with V = 0. In other words, this corresponds to (e &eb) a particle traveling freely through (e&eb) empty space. Time-dependent** |
| **Schrödinger equation (general)** |
| **(16) i ℏ ∂ ∂ t Ψ ( r , t ) = H ^ Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t) where i is the imaginary unit, ħ is the Planck constant divided by 2π, the symbol ∂/∂t indicates a partial derivative with respect to time t, Ψ (the Greek letter psi) is the wave function of (e) the quantum system, and Ĥ is the Hamiltonian operator (which characterizes the total energy of any given wave function and takes different forms depending on the situation).** |
| **(17) Each of these three rows is a wave function which satisfies (eb) the time-dependent Schrödinger equation for a harmonic oscillator. Left: The real part (blue) and imaginary part (red) of the wave function. Right:** |
| **(18) The probability distribution of finding the particle with (e&eb) this wave function at a given position.** |
| **(19) The top two rows are examples of stationary states, which correspond to (e&eb) standing waves.** |
| **(20) The bottom row is an example of a state which is (=) not a stationary state.** |
| **(21) The right column illustrates (eb) why stationary states are (=) called "stationary".** |

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| **(22) The most famous example is the non-relativistic Schrödinger equation for (e) a single particle moving in an electric field (but not a magnetic field; see the Pauli equation): [6]** |
| **Time-dependent Schrödinger equation** |
| **(Single non-relativistic particle)** |
| **(23) i ℏ ∂ ∂ t Ψ ( r , t ) = [ − ℏ 2 2 μ ∇ 2 + V ( r , t ) ] Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)** |
| **(24) Here μ is (=) the particle's "reduced mass", V is its potential energy, ∇2 is the Laplacian (a differential operator), and Ψ is (=) the wave function (more precisely, in this context, it is called the (=) "position-space wave function").** |
| **(25) In plain language, it means "total energy equals (=) kinetic energy plus potential energy", but the terms take unfamiliar forms for (e) reasons explained below.** |
| **(26) Given the particular differential operators involved, this is (=) a linear partial differential equation.** |
| **(27) It is also a diffusion equation, but unlike the heat equation, this one is (=) also a wave equation given (e&eb) the imaginary unit present in the transient term.** |
| **(28) The term "Schrödinger equation" can refer to (e) both the general equation (first box above), or the specific nonrelativistic version (second box above and variations thereof).** |
| **(29) The general equation is indeed quite general, used throughout quantum mechanics, for everything from the Dirac equation to quantum field theory, by plugging in various complicated expressions for the Hamiltonian. The specific nonrelativistic version is a simplified approximation to reality, which is quite accurate in many situations, but very inaccurate in others (see relativistic quantum mechanics and relativistic quantum field theory)** |
| **(30) To apply the Schrödinger equation, the Hamiltonian operator is set up for the system, accounting for the kinetic and potential energy of the particles constituting the system, then inserted into the Schrödinger equation. The resulting partial differential equation is solved for the wave function, which contains information about the system.** |
| **Time-independent equation** |

**(31) The time-dependent Schrödinger equation described above predicts that wave functions can form standing waves, called stationary states (also called "orbitals", as in atomic orbitals or molecular orbitals). These states are important in their own right, and if the stationary states are classified and understood, then it becomes easier to solve the time-dependent Schrödinger equation for any state. Stationary states can also be described by a simpler form of the Schrödinger equation, the time-independent Schrödinger equation. (This is only used when the Hamiltonian itself is not dependent on time explicitly. However, even in this case the total wave function still has a time dependency.)**

**Time-independent Schrödinger equation (general)**

**(32) E Ψ = H ^ ⊠ Ψ {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi } {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi }**

**In words, the equation states:**

**(33) When the Hamiltonian operator acts on a certain wave function Ψ, and the result is proportional to the same wave function Ψ, then Ψ is a stationary state, and the proportionality constant, E, is the energy of the state Ψ.**

**(34) The time-independent Schrödinger equation is discussed further below. In linear algebra terminology, this equation is an eigenvalue equation.**

**(35) As before, the most famous manifestation is the non-relativistic Schrödinger equation for a single particle moving in an electric field (but not a magnetic field):**

**Time-independent Schrödinger equation (single non-relativistic particle)**

**(36) E Ψ ( r ) = [ − ℏ 2 2 μ ∇ 2 + V ( r ) ] Ψ ( r ) {\displaystyle E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )} E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )**

**Gravitational waves: (Wikipedia)**

**(37) In theory, the loss of energy through gravitational radiation could eventually drop the Earth into the Sun. However, the total energy of the Earth orbiting the Sun (kinetic energy + gravitational potential energy) is about 1.14×1036 joules of which only 200 Watts (joules per second) is lost through gravitational radiation, leading to a decay in the orbit by about 1×10−15 meters per day or roughly the diameter of a proton. At this rate, it would take the Earth approximately 1×1013 times more than the current age of the Universe to spiral onto the Sun. This estimate overlooks the decrease in r over time, but the majority of the time the bodies are far apart and only radiating slowly, so the difference is unimportant in this example.**

**More generally, the rate of orbital decay can be approximated by[39]**

**(38)** $\frac{dr}{dt} = -\frac{64}{5}\frac{G^3}{c^5}\frac{(m_1 m_2)(m_1 + m_2)}{r^3}$ , {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,} {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,}

**where r is the separation between the bodies, t time, G Newton's constant, c the speed of light, and m1 and m2 the masses of the bodies. This leads to an expected time to merger of [39]**

$t = \frac{5}{256}\frac{c^5}{G^3}\frac{r^4}{(m_1 m_2)(m_1 + m_2)}$. {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.} {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.}

**(39) For example, a pair of solar mass neutron stars in a circular orbit at a separation of 1.89×108 m (189,000 km) has an orbital period of 1,000 seconds, and an expected lifetime of 1.30×1013 seconds or about 414,000 years. Such a system could be observed by LISA if it were not too far away. A far greater number of white dwarf binaries exist with orbital periods in this range. White dwarf binaries have masses in the order of the Sun, and diameters in the order of the Earth. They cannot get much closer together than 10,000 km before they will merge and explode in a supernova which would also end the emission of gravitational waves. Until then, their gravitational radiation would be comparable to that of a neutron star binary.**

**(40) By substituting value of dr/dt in Schrödinger's wave equation we get the quantum wave equation for gravitational waves.**

In order to have (H3) and (H4) hold, $m, \mu_0$, K, b, δ, and $\omega$ need to satisfy

$$\frac{\omega_1}{(m + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 + b_1} + \sigma)^2} + \frac{2\omega_1}{([m_1]_1 + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 + b_1} + \sigma_1)^3} < 1$$

And

$$\frac{1}{(m_1 + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 - b_1})^2} + \frac{2\omega_1}{(m_1 + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 - b_1})^3} > 1$$

When the inequalities (33) and (34) hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$(m_1 + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 - b_1} + (m_1 + [\mu_0]_1 + \frac{K_1 b_1}{\delta_1 - b_1} + \sigma_1)\, Q_1^*)^3$

$= m_1 + [\mu_0]_1 + 3 + \frac{K_1 b_1}{\delta_1 - b_1}(1 + Q^*)(1 + \omega_1 Q_1^*$

Now

$$\left[a_1^2\right] - 4\left[a_2\right] e = \left(\frac{2K_1 b_1}{\delta_1 - b_1}\right)^2 + \frac{4K_1 b_1 \sigma_1 Q_1^*}{(\delta_1 - b_1)(1 + Q_1^*)} + \frac{\sigma_1^2 Q_1^{*2}}{(1 + Q_1^*)^2}$$

$$\frac{4\, b_1 \sigma_1 (\delta_1 - b) Q^*}{\delta_1 (1 + Q_1^*)^2} - \frac{4\, (\delta_1 - b_1)^2 (1 - \omega_1) Q_1^*}{\delta_1 (1 + \omega_1 Q_1^*)(1 + Q_1^*)} P^*(0)$$

And

$\Delta = \lim_{0 < a < \infty} |[h_*]_1\, a e_1^{mu}\ [\mu_0]_1 - \frac{a_1}{a_1 + 1}$

Hence if $[a_1^2 > 4(a_2)_1 e$ and

$$[a_1] = \frac{2K_1 b_1}{\delta_1 - b_1} + \frac{\sigma_1 Q_1^*}{1 + Q_1^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_1]_1 = [f_{1*}]_1 + [f_{2*}]_1 + [f_{1*}] P_1^*$$

$$[a_2]_1 = \left[f^{1}_{2*}\right] P^{*}_{1} - \frac{h^{1}_{*}}{h_{*}} P^{*}(0)$$

$$e = [g^{1}_{*}]_1 Q^{*}_{1}$$

$$\eta_1(t) = \left[h_{*}\right]_1 \beta_1 \left(\theta_1(t)\right) - \mu_1(\theta(t)),$$

The system(31)can be written as the following system:

$$[M(t)]\left[a_1\right]M_1 - [a_2]_1 W_1 - [n(t)]M_1$$

$$w_1(t) = eM_1$$

Write

$$A = \begin{pmatrix} -[a_1]_1 - [a_2]_1 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)= \begin{pmatrix} [\eta(t)]_1 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t) = \begin{pmatrix} M_1 \\ W_1 \end{pmatrix}$$

The solution of (32)is

$$X_1(t) = e^{A,t}[X]_1(0) + \int_0 e^{At}$$

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# Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

## (Module Three)

**Abstract: Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-435) 1990. We demonstrate its applicability cosmology/ quantum computation economics et al. As stated in the foregoing we shall apply model to various cosmological problem and show the relevance to the same. . To exemplify model can be applied to the following situational significations: Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an**

**outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 Masaru Shibata and Takashi Nakamura** perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. **Authors** use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society **Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995**

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(g_3)Q_3$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_3(0) = 0, \quad g_3(P_3) > 0, \quad P_3 \text{ 0}, \qquad \lim_{P_3 \to \infty} g_3(P_3) = (K_3)_3 \text{ ∞}. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_3 + (\rho_a)_3 + D_3(a_3, P_3, Q_3)\rho_3 + 0, \qquad 0 < a_3 \text{ ∞},$$

$$\rho_3(t,0) = \int_0^\infty n_3(a_3, P_3, Q_3)\rho_3(t_3, a_3)da_3$$

Where $D_3$ is the mortality function and $n_3$ is the fertility function.

Put $_3$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_3}{dt}\right]_3 = x_3\left[\frac{dQ_2}{dt}\right]_3 - x_2\left[\frac{dQ_3}{dt}\right]_3$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_3 - x_4\left[\frac{dQ_2}{dt}\right]_3$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_3 - x_6\left[\frac{dQ_3}{dt}\right]_3$$

Where $x_3, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_3$ and $g_3(P_3)Q_3$

Now when you add the term $(b_3Q_3)$ with $g_3(P_3)(Q_3)$ we are infect see is that $(b_3Q_3)$ is credited incrementally part of $g_3(P_3)(Q_3)$ or instant nearly, when subtracted from $\frac{dQ_3}{dt}$ , it means that $g_3(P_3)(Q_3)$ is slowly eaten by $\frac{dQ_3}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_3 = \mu_3(a_3) + (f_3)_3(P_3)_3 + (f_2)_3(Q)_3 \quad (2)$$

$$n_3 = \beta_3(a_3)h_3(Q_3), \quad (3)$$

Where $\mu_3$the age- dependent debit rate whit is $\mu_3(a_3)$ 0 $\forall a > 0$; $(f_3)_3$ is the density- dependent death rate with $f_3(0) = 0$, $(f_3^3)_3(P_3)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_3)_3(P_3) = (k_3)_3$ ; $f_2$ is the functional response of

creditors to prey with $[f_2]_3(0) = 0, (f_2^3)_3(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_3\ (Q_3) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 3, h^3(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_3 + (\rho_a)_3 = -\mu_3(a_3) + (f_3)_3 P_3) + \big(f_2)_3(Q_3)\big)(\rho_3)_3 \quad (4)$$

$$\rho_3(t_3, 0) = h_3(Q_3)\int_0^\infty \beta_3\,(a_3)\rho_3(t, a_3)da, \quad (5)$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(P_3)Q_3 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_3$ and $Q_3$ of (4-6) satisfy

$$0 \ \rho_3(t, a_3) \ \rho_3^*(t, a_3),$$

$$0 \ Q_3(t) \ Q_3^*(t),$$

Where $\rho_3^*(t, a_3)$ and $Q_3^*(t)$, are solutions of

$$(\rho_t^*)_3 + (\rho_a^*)_3 = -(\mu_3(a_3) + (f_3)_3(P_3^*) + (f_2)_3 Q_3^*))\,\rho_3^*, \quad (7)$$

$$\rho_3^*\,(t, 0) = \int_0^\infty \beta_3(a_3)\,\rho_3^*\,(t, a_3)da_3 \quad (8)$$

$$\frac{dQ_3^*}{dt} = -\,b_3\,Q_3^* + g_3(P_3^*)Q_3^*, \quad (9)$$

With $P_3^*(t) = \int_0^\infty \rho_3^*\ (t, a_3)da_3$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_3^*(t)}{dt} = P_3^*(t)\big(r_3 - (f_3)(P_3^*) - (f_2)_3(Q^*)\big), \quad (10)$$

$$\frac{dQ_3^*(t)}{dt} = Q_3^*(t)(-b_3 + g_3(P_3^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_3(a)e^{-\int_0^a \mu_3(\tau)dt}e^{-ra}\ da_3 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_3^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(t)dt}\ da_3 > 1, \quad (H1)$$

Which guarantees $r_3 > 0$.

*Remark.* If $\beta_3(a) = [\beta_0]_3 a_3 e^{-\alpha_3 a_3}, \mu_3(a_3) \equiv \mu_3$ then the solution of (12) is

$$r_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_3]_3(P_3) = 0, [f_2](Q_3) = \lambda_3 Q_3\ h_3(Q_3) = 1, \quad and\ \ g_3(P_3) = c_3 P_3$$

$$[\rho_3]_3 + [\rho_a]_3 = -(\mu_3 + \lambda_3 Q_3)\rho_3$$

$$\rho_3(t,0) = \int_0^{\infty} [\beta_0]_3 a_3 e^{-\alpha_3 a_3} \rho_3(t,a_3)\, da_3$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + c_3 P_3 Q_3$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_3}{dt} = \omega_3 P_3 - \lambda_3 P_3 Q_3$$

$$\frac{dQ_3}{dt} = b_3 Q_3 - c_3 P_3 Q_3$$

Where $\omega_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$

## THE EQUILIBRIUM $(\rho_0)_3(a_3), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_3(a_3), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_3 + (\rho_a)_3 = -\big(\mu_3(a_3) + (f_3)_3(P_3)\big)\rho,$$

$$\rho_3(t,0)$$

$$= \int_0^\infty \beta_3\,(a_3)\rho_3(t,a_3)da_3 \qquad (13)$$

The solution of (13) has the form of

$$\rho_3(t,a_3) = Ce_3^{r(t-a_3)} - \int_0^a (\tau)d\tau - \int_0^3 (f_3)_3\big(P_3(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_3(t) = C_3 e^{r_3 t - \int_0^t (f_3)_3(P_3(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_3^{-3}(r)$ if $k_3 > r$. This is equivalent to

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{(k_3)_3 a_3} da_3 < 1.$$

In fact, if we define

$$R_3(s_3) = \int_0^\infty \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau)d\tau} e^{-s_3 a_3}\, da_3$$

$R(s)$ is a decreasing function. $R_3(r_3) = 1$ And $R_3(k_3)_3 < 1$ imply $r_3 < (k_3)_3$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_3(a_3), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{k_3 a_3} da_3$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_3 \equiv \int_0^\infty (\rho_0)_{33}(a_3)da_3$ *satisfies*

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{-(f_3)_3(P_0)a_3} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_3(a_3), 0)$

Let $u_3(t, a_3) = \rho(t, a_3) - \rho_0(a_3)$ Then

$$N_3(t) \equiv \int_0^{\infty} u_3(t, a_3) da_3 = \int_0^{\infty} \rho_3(t, a_3) da_3 - \int_0^{\infty} (\rho_0)_3(a_3) da = P_3(t) - P_0)_3$$

And

$$(u_3)_3 + (u_a)_3 = (\rho_t)_3 + (\rho_a)_3 - (\rho_0^3)_3(a_3)$$

$$= -\mu_3(a_3) + (f_3)_3(P_3) + (f_2)_3(Q_3))(u_3 + (\rho_0)_3) + \mu_3(a_3)(\rho_0)_3 + (f_3)_3(P_0)_3(\rho_0)_3$$

$$= -\mu_3(a_3)u_3 - (f_3)_3(P_3) + (f_2)_3(Q_3))u_3 - (\rho_0)_3 (f_3)_3(P_3) - (f_3)_3(P_0)_3$$

$$+ (f_2)_3(Q_3)), \qquad (15)$$

Where $(\rho_0^3)_3(a_3)$ denotes $d(\rho_0)_3(a_3)/da_3$

Expansion $(f_3)_3(P_3)$ and $(f_2)_3(Q_3)$ at $(P_0)_3$ and 0 yield respectively

$$(f_3)_3(P_3) = (f_3)_3((P_0)_3 + N_3)$$

$$= (f_3)_3(P_0)_3 + (f_3^3)_3(P_0)_3 N_3 + o_3(N_3), \qquad (16)$$

$$(f_2)_3(Q_3)$$

$$= (f_2^3)_3(0)Q + o_3(Q_3) \qquad (17)$$

Write $f_3(P_0) = f_{30}, f_3'(P_0) = f_{30}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_3)_3 + (u_a)_3 \approx -(\mu_3(a_3) + (f_{30})_3)u - (f_{30}^3)_3(\rho_0)_3(N_3) - (f_{20}^3)_3(\rho_0)_3(Q_3)$$

$$u_3(t, 0) = \rho_3(t, 0)(\rho_0)_3(0)$$

$$= h_3(Q_3)\int_0^{\infty} \beta_3(a_3)(u_3(t,a) + (\rho_0)_3(a_3)da_3 - (\rho_0)_3\,(0)$$

$$= h_3(Q_3)\int_0^{\infty} \beta_3(a_3)(u_3(t,a_3)da_3 + (h_3(Q_3) - 3)(\rho_0)_3\,(0) \qquad (18)$$

Expanding

$$h_3(Q_3) = 1 + \ (h^3)_3(0)Q_3 + 0(Q_3) \equiv 1 + h_0^{'}Q_3 + 0Q_3$$

And substituting it into (38), we arrive at

$$u_3(t,0) \approx \int_0^{\infty} \beta_3\ (a_3)u_3(t,a_3)da_3 + (h_{a_3}^{'})(\rho_0)_3Q_3$$

The linearization is the following :

$$[u_3] + [u_3] = -(\mu_3(a_3) + [t_{30}]_3u_3 - [\alpha_3]_3(a_3)N_3 - [\alpha_2]_3(a_3)Q_3$$

$$u_3(t,0) = \int_0^{\infty} \beta_3(a_3)u_3(t,a_3)da_3 + [\xi_3]Q_3 \qquad (19)$$

$$\frac{dQ_3}{dt} = -b_3Q_3$$

where

$$[\alpha_3]_3(a_3) = [t_3^3a]_3[p,o]_3(a_3)d2(a_3) = [f_{20}^3]_3(\rho\ o\ )_3(a_3)\xi = [\mathrm{h\ o}^3]_3[\rho\mathrm{o}]_3(\mathrm{o}) \text{ and } \mathrm{b}_3 = \mathrm{b}_3 - g_3(P_3)$$

If $g_3(P_0) > b$, clearly, $(\rho_0)_3(a_3), 0)$ is unstable and the predator population is persistent. Assume $g_3(P_0)_3 < b$.

Then, $Q_3(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_3 + u_a = -(\mu_3(a_3) + (f_{30})_3)u_3 - (\alpha_3)_3(a_3)N_3$$

$$= -(\mu_3(a_3) + f_{30})u_3 - \alpha_3(a_3)\int_0^\infty u_3(t, a_3)da_3$$

$$u_3(t, 0) = \int_0^\infty \beta_3\,(a_3)u_3(t, a_3)da_3$$

Using the method developed by Cushing [2, 3], we suppose

$$u_3(t, a_3) = \tilde{u}_3(a_3)^{e^{\gamma(t-a_3)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_3}\,da_3$. Then,

A solution of the first differential equation for $\tilde{u}(a_3)$ gives,

$$\tilde{u}_3(a_3) = e^{-\int_0^a \mu_3(\tau_3)d\tau - (f_{30})_3 a}\left(\tilde{u}_3(0) - w\int_0^a e^{\int_0^\tau \mu_3(s_3)ds + (f_{30})}\tau_3\,(\alpha_3)_3(\tau_3)e^{\gamma_3\tau}d\tau\right).$$

Substituting the expression (13) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\right)\tilde{u}_3(0)$$

$$+w\int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\int_0^a (\alpha_3)_3(\tau_3)e^{-\gamma_3\tau_3}e^{(f_{30})_3\tau_3 + \int_0^\tau \mu_3(s_3)ds}d\tau_3\,da_3$$

$$= 0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(\tau_3)d\tau} da\right)(\tilde{u}_3(0))_3$$

$$+\left(-3-\int_0^\infty e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(\tau_3)d\tau}\int_0^a e^{(f_{30})_3\tau_3+\int_0^\tau \mu_3(s_3)ds+y\tau}\alpha_3(\tau)d\tau\, da\right)w=0. \quad (23)$$

Define $Y(\gamma,a)=e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(t_3)d\tau}$. Then

$$\left(1-\int_0^\infty \beta_3\,(a_3)Y_3(\gamma_3,a_3)da_3\right)u_3(0)$$

$$+w\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)\int_0^a(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\, da_3=0, \quad (24)$$

$$\left(\int_0^\infty Y_3(\gamma_3,a_3)da_3\right)\tilde{u}_3(0)$$

$$+\left(-1-\int_0^\infty Y_3(\gamma_3,a_3)\int_0^a(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\, da_3\right)w=0.$$

Equqtion (20) has a solution of the form of $u_3(t,a_3)=u_3(a_3)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_3(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_3(\gamma_3,a_3)=\int_0^a(\rho_0)_3(\tau_3)Y^{-3}\,(\gamma_3,a_3)d\tau_3$. We obtain the characteristic equation

$$f_{30_3}^l\left(\int_0^\infty Y_3(\gamma_3,a_3)da_3\right)\left(\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$+\left(1-\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)da\right)\left(1+(f_{30}^l)_3\int_0^\infty Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_3(a_3)=(\rho_0)_3(0)e^{-\int_0^a \mu_3(\tau_3)d\tau_3-(f_{30})_3a_3}$$

$$I_3(\gamma_3,a_3)=\int_0^a (\rho_0)_3(\tau_3)e^{\gamma_3\tau_3+(f_{30})_3\tau+\int_0^\tau \mu_3(s_3)ds_3}d\tau$$

$$=\frac{(\rho_0)_3(0)}{\gamma_3}(e^{\gamma_3a_3}-1)$$

Let

$$E_3(a_3)=e^{-(f_{30})_3a_3-\int_0^a \mu_3(\tau_3)d\tau_3}$$

Then $[Y_3]_3[\gamma_3]_3,a_3)=E_3(a_3)e^{-\gamma_3a_3}$. by a simple calculation,

$$\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)I(\gamma_3,a_3)da=\frac{[\rho_0]_3(0)}{\gamma}\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3a_3}\ da\right);$$

$$\int_0^\infty Y_3(\gamma_3,a_3)I(\gamma_3,a_3)da=\frac{[\rho_0]_3(0)}{\gamma_3}\left(\int_0^\infty E_3(a_3)e^{-\gamma_3a_3}\ da_3\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_3(0)}{\gamma_3}[f_{30}^3]\int_0^\infty E_3(a_3)e^{-\gamma_3a_3}da_3\left(\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3a_3}da_3\right)$$

$$+\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right).$$

$$\left(1+\frac{(\rho_0)_3(0)}{\gamma_3}(f_{30}^{'})_3\left(\int_0^\infty E_3(a_3)da_3-\int_0^\infty E_3(a_3)e^{-\gamma_3 a_3}da_3\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right)(1+\frac{(f_{30}^{'})_3}{\gamma_3}(\rho_0)_3(0)\int_0^\infty E_3(a_3)da_3)=0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_3(a_3)Y_3(\gamma_3 a_3)da_3 = 1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3 \neq 0,$$

Since $\int_0^\infty \beta_3(a_3)Y_3(\gamma_3 a_3)\int_0^a \alpha_3(\tau_3)Y^{-3}(\gamma_3,\tau_3)d\tau_3 da_3 > 0$. Therefore, the characteristic solution is

$$\gamma_3 = -(f_{30}^{'})_3(\rho_0)_3\int_0^\infty E_3(a_3)da_3.$$

Since $f_3'(P_0)_3 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_3), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_3), 0)$ which is related to the extinction of debit summation types

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_3 + (\mu_3(a_3) + [f_3]_3(P_3^*) + [f_2]_3(Q_3^*))P_3^* = 0, \quad (26)$$

$$p_3^*(0) = h_3(Q_3^*)\int_0^\infty \beta_3(a_3)p_3^*(a_J)p_3^*(a_3)da_3 \quad (27)$$

$$g_3(P_3^*) - b_3 = 0, \quad (28)$$

Where $P_3^* = \int_0^\infty p_3^*(a_3)\, da_3$

Equation (28) yields $P^* = g^{-3}(b)$,

From (26),

$$p_3^*(a) = p_3^*(0)e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}$$

Substituting into (27), we need to solve

$$\left(1 - h_3(Q_3^*)\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}da_3\right)p_3^*(0) = 0$$

For $Q_3^*$

Fix $p_3^*(0)$ and $P_3^*$ Let

$$H_3(z_3) = 1 - h_3(z_3)\int_0^{\infty} \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3} da_{,3}$$

And assume that

$$[k_3]_3 \int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(g_3^{-3}(b_3))a_3 - [k_2]_3 a_3}\, da_3 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_3(z_3)] > 0.$$

Clearly,

$$\frac{dH_3}{dz_3} = h_3(z)([f_2^{'}]_3(z_3))\int_0^{\infty} a_3\beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - f_3(P_3^*)a_3 - f_2(z_3)a_3} da_{\square}$$

$-h^{'}(z_3)\int_0^{\infty} \beta_3(a_3)\, e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3}\ da_3$

By the assumptions on $\left[f_2^3\right]_3 and\ [h^3]^3 dH_3/dz > 0$. Hence, if

$H_3(0) = 1 - \int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3} da_3 < 0,$ (29)

There exists a unique positive solution $Q^*$

Now, we claim that $H_3(0) < 0$ follows if $g_3^{-}(b_3) < [f_3^{-3}]_3(r_3)$, where $r_3$ is the solution of (12).

Indeed, since

$$R_3(s_3) = \int_0^{\infty} \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 e^{-y\frac{a}{3}} da_3}$$

Is a decreasing function of $s_3$ when $P_3^* = [g^{-3}]_3(b_3) < [f_3^{-3}]_3(r_3)$, that is $[f_3]_3(P_3^*) < r$, $R\big(f_3(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_3(g^{-3}(b))) > 1, i.e.,$

$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3\left(g_3^{-3}(b)\right)a_3} > 1.$ (H4)

Then the system (4)-(6) has a unique positive equilibrium $(P_3^*(a_3), Q_3^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_3^*(a_3), Q_3^*)$ is more complicated. We again consider perturbations $m_3(t_3 a_3) = p_3(t_3, a_3) - p_3^*(a_3)$ and $W_3(t) = Q_3(t) - Q^*$.Then

$M_3(t) \equiv \int_0^{\infty} m_3(t, a_3) da_3 = P_3(t) - P_3^*,$

$$[m_3]_3 + [m_a]_3 = [p_3]_3 + [p_a]_3 - \left[p^*\right]_3(a_3) \approx -(\mu_3(a_3) + [f_{3*}]_3 + [f_{2*}]_3)m_3$$

$$-[f_{3*}^3]_3 p_3^*(a_3)M - [f_{2*}^3]_3 p_3^*(\mathrm{a})W_3,$$

$$m_3(t,0)p_3(t,0) - p_3^*(0) \approx h_* \int_0^{\infty} \beta_3(a_3) m_3(t, a_3)\, da + \frac{[h_*^3]_3}{[h_*]_3}\, p_3^*(0)W_3$$

$$\frac{dW_3}{dt} = \frac{dQ_3}{dt} \approx [g_*^3]_3 Q_3^* M_{,3}$$

Where $[f_{3*}]_3$ denoting $[f_3(P^*)]_3$, $[f_{3*}]_3$ denoting $[f_2(P^*)]_3$, $[f'_{3*}]_3$ denoting $[f'_3(P^*)]_3[f'_{2*}]_3$ denoting $[f_2]_3$ $(Q_3^*,\quad h_3^*$ denoting $h'_*$ denoting $h'_3$ $(Q_3^*)$, and $[g'_*]_3$

Denoting $[g']_3(P_3^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_3]_3 + [m_a]_3 = -(\mu_3(a_3) + [f_{3*}]_3 + [f_{2*}]_3)m_3 - [f'_{3*}]_3[p_3^*](a_3)M_3 - [f'_{2*}]_3 p_3^*(a_3)W_3$$

$$m_3(t,0) = [h_*]_3 \int_0^\infty \beta_3(a_3)m_3(t,a_3)da_3 + \frac{[h_*^3]_3}{[h^*]_3}[p^*]_3(0)\,W_{,3} \qquad (30)$$

$$\frac{dW_3}{dt} = [g_*^3]_3[Q_3^*]M_3$$

Integrating $[m_3]_3(t,a_3)$ with respect to a from 0 to ∞, we have

$$\frac{dM_3(t)}{dt} = \int_0^\infty [mt]_3(t,a_3)da_3$$

$$= -\int_0^\infty [m_a]_3\,(t,a_3)da_3 - \int_0^\infty \mu(a_3)m(t,a_3)da_3$$

$$-\left([f_{3*}]_3 + [f_{2*}]_3 + [f'_{3*}]_3\int_0^\infty p_3^*(a_3)da_3\right)M_3 - [f_{2*}^3]_3\int_0^\infty p_3^*(a_3)da_3 W_3$$

$$= m_3(t,0) - \int_0^\infty \mu_3(a_3)\,m_3(t,a_3)da_3$$

$$-\left([f_{3*}]_3 + [f_{2*}]_3 + [f'_{3*}]_3 P_3^*\right)M_3 - [f'_{2*}]_3 P_3^* W_{,3}$$

From (30), we arrive at the equations

$$\frac{dW_3(t)}{dt} = \int_0^{\infty} ([h_*]_3\beta_3(a_3) - \mu_3(a_3))m_3(t,a_3)da_3$$

$$-\left([f_{3*}] + [f_{2*}]_3 + \left[f_{3*}^{'}\right]_3 P_3^*\right) M_3 + \left(\frac{[h_*^{'}]_3}{[h_*]^3}\, P_3^*(0) - f_{2*}^{'} P_3^*\right) W_3,$$

(33)

$$\frac{dW_3(t)}{dt}\left[g^{'}{}_*\right]_3 Q_3^*\, M_{.3}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{"*}]_3\beta_3(a_3) - \mu_3(a_3))m_3(t,a_3)da_3 = ([h_*]_3\beta_3(\theta_3) - \mu_3(\theta_3))M_{,3}$$

Where $0< \theta_3 = \theta_3(t) < \infty$.

And if $a_3^2\,[a_3^2]_3<4\,[a_2]_3<4\,a_2 e$,

Re $\lambda_3\left[\lambda_3\right]_3 = \text{Re}\lambda_2\left[\lambda_2\right]_3 = -\frac{3}{2}a_3[a2\ ]_3 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_*[h_*]_3\beta_3(a_3) - \mu_3(a_3)\right|$. Then$|n_3(t)|$ ⊠⊿∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_3^2 <4\,a_2 e$ and $[a_3] > 2\Delta$,or $[a_3^2]$ ⊠4$[a_2]e$ and $(a_3 - \sqrt{(a_3^2)_3 - 4(a_2)_3 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+3}\right)$ and $\beta_3(a_3) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3) + d\tau_3 da} = \frac{3}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_3 + [\mu_0]$ 0.5. Assume$[f_3](x_3)= K_3x_3$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_3(x_3)= [0x_3]/(3 + x_3)$; the functional responses $[[f_2](x) = [0]\, x_3/(1 + x_3)$and $h_3(x_3)=(1+\omega_3x_3)/(1 + x_3)$(See Holling [7] and Rashevsky [13])

Equation (34),now has the form of

$$\frac{m_3 + [\mu_0] + [K][P_0])}{[m_3]_3 + (\mu_0)_3 + [K][P_0])_3^3\ [K]_3[P_0]_3^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(3 + P_0) < b_3$i e $[P_0] < b/(\delta_3 - b_3)$then the equilibrium$(P_0)$ $[P_0]_3(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_3$

$[f_{3*}] = K_3\ P^*\ [f_{2*}] = \frac{\sigma Q_3^*}{3+Q_3^*}$ $[h_*] = \frac{3+\omega_3 Q_3^*}{3+Q_3^*}$

$[f_{3*}] = K_3\ [f_{2*}^3] = \frac{\sigma_3}{(3+Q_3^*)^2}[\ g_*^3] = \frac{\sigma_3}{(3+P_3^*)^2}$ $[h_*^3] = \frac{\omega_3-3}{(3+Q_3^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0, K, b, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_3}{(m+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma)^2}+\frac{2\omega_3}{([m_3]_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma_3)^3}<1$$

And

$$\frac{1}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^2}+\frac{2\omega_3}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}+(m_3+[\mu_0]_3+\frac{K_3b_3}{\delta_3-b_3}+\sigma_3)\,Q_3^*)^3$$

$$=m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}(1+Q^*)(1+\omega_3Q_3^*$$

Now

$$\left[a_3^2\right]-4\left[a_2\right]e=\left(\frac{2K_3b_3}{\delta_3-b_3}\right)^2+\frac{4K_3b_3\sigma_3\,Q_3^*}{(\delta_3-b_3)(1+Q_3^*)}+\frac{\sigma_3^2\,Q_3^{*2}}{(1+Q_3^*)^2}$$

$$\frac{4\,b_3\sigma_3\,(\delta_3-b)Q^*}{\delta_3\,(1+Q_3^*)^2}-\frac{4\,(\delta_3-b_3)^2(1-\omega_3)Q_3^*}{\delta_3(1+\omega_3Q_3^*)(1+Q_3^*)}P^*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_3\,ae_3^{mu}\;\;[\mu_0]_3-\frac{a_3}{a_3+3}$$

Hence if $[a_3^2 > 4(a_2)_3 e$ and

$$[a_3] = \frac{2K_3 b_3}{\delta_3 - b_3} + \frac{\sigma_3 Q_3^*}{1 + Q_3^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_3]_3 = [f_{3*}]_3 + [f_{2*}]_3 + [f_{3*}] P_3^*$$

$$[a_2]_3 = \left[f_{2*}^3\right] P_3^* - \frac{h_*^3}{h_*} P^*(0)$$

$$e = [g_*^3]_3 Q_3^*$$

$$\eta_3(t) = \left[h_*\right]_3 \beta_3 \left(\theta_3(t)\right) - \mu_3(\theta(t)),$$

The system(33)can be written as the following system:

$$[M(t)]\left[a_3\right] M_3 - [a_2]_3 W_3 - [n(t)] M_3$$

$$w_3(t) = eM_3$$

Write

$$A = \begin{pmatrix} -[a_3]_3 - [a_2]_3 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_3 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_3 \\ W_3 \end{pmatrix}$$

The solution of (32)is

$$X_3(t) = e^{A,t}[X]_3(0) + \int_0 e^{At}$$

**This equation is useful in following type of problems . For actual model kindly see various papersThere is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant Λ, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive Λ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the Λ>0 generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible.**

**Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of (e) GW signals from BBHs with component spins $\chi_{1,2} \in [-0.99, 0.99]$ aligned (e&eb) with the orbital angular momentum, component masses $m_{1,2} \in [2,48] M_\odot$, and (e&eb) total mass $M_{total} \leq 50\ M_\odot$. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins >+0.6 and component masses $\gtrsim 20\ M_\odot$. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201 STOP Advanced LIGO may be the first experiment to (e) detect gravitational waves. Through superradiance of stellar black holes, it may also be (=) the first experiment to discover the QCD axion with (e&eb) decay constant above the GUT scale. When an axion's Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e) the black hole, forming a "gravitational atom." Through the superradiance process, the number of axions occupying (e) the bound levels grows exponentially, extracting (e) energy and angular momentum from the black hole. Axions**

**transitioning between levels of the gravitational atom and (e&eb) axions annihilating to (e) gravitons can produce (eb) observable gravitational wave signals. The signals are long-lasting, monochromatic, and can be distinguished from (e) ordinary astrophysical sources. Authors estimate up to O (1) transition events at aLIGO for (e) an axion between 10^-11 and (e&eb) 10^-10 eV and up to 10^4 annihilation events for (e) an axion between 10^-13 and 10^-11 eV. In the event of a null search, aLIGO can constrain (e) the axion mass for (e) a range of rapidly spinning black hole formation rates. Axion annihilations are also promising for (e) much lighter masses at future lower-frequency gravitational wave observatories; the rates have (e) large uncertainties, dominated by (e) supermassive black hole spin distributions. Projections for aLIGO are (=) robust against perturbations from (e) the black hole environment and account for (e) updated exclusion on (e&eb) the QCD axion of 6*10^-13 eV < ma < 2*10^-11 eV suggested by (e) stellar black hole spin measurements. Subjects: High Energy Physics - Phenomenology (hep-ph); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 91, 084011 (2015) DOI: 10.1103/PhysRevD.91.084011Cite as:arXiv:1411.2263 [hep-ph] (or arXiv:1411.2263v3 [hep-ph] for this version) Discovering the QCD Axion with Black Holes and Gravitational Waves Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang STOP Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric**

**charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by (e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave**

**signal. The signal sweeps upwards in frequency from 35 to 250 Hz with (e&eb) a peak gravitational-wave strain of (e) 1.0×10−21. It matches the waveform predicted by general relativity for (e) the inspiral and merger of a pair of black holes and the ringdown of the resulting (eb) single black hole. The signal was observed with a matched-filter signal-to-noise ratio of (e) 24 and a false alarm rate estimated to be (=) less than 1 event per 203 000 years, equivalent to (=) a significance greater than 5.1σ. The source lies at (eb) a luminosity distance of 410+160−180 Mpc corresponding to (e&eb) a redshift z=0.09+0.03−0.04. In the source frame, the initial black hole masses are (=) 36+5−4M⊙ and 29+4−4M⊙, and the final black hole mass is (=) 62+4−4M⊙, with (e&eb) 3.0+0.5−0.5M⊙c2 radiated in (eb) gravitational waves. All uncertainties define (eb) 90% credible intervals. These observations demonstrate (eb) the existence of binary stellar-mass black hole systems. Observation of Gravitational Waves from a Binary Black Hole Merger B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 061102 – Published 11 February 2016 This is the first direct detection of gravitational waves and (e&eb) the first observation of a binary black hole merger. Received 21 January 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.061102 STOP This is the second in a series of three papers on (e&eb) systematic analysis of rank 1 Coulomb branch geometries of (e) four dimensional N=2 SCFTs. In the first paper authors developed a strategy for (e) classifying physical rank-1 CB geometries of N=2 SCFTs. They show how to carry out this strategy computationally to construct (eb) the Seiberg-Witten curves and (e&eb) one-forms for all the rank-1 SCFTs. Explicit expressions are given for all cases, with the exception of the Nf=4 SU (2) gauge theory and (e&eb) the En SCFTs which were previously constructed. Classification includes all known rank-1 theories plus (+) a new one with (e&eb) an abelian flavor group, plus (+) nine additional theories whose existence is more speculative. Four of those, reported in first paper, depend on (=) the assumption of new frozen rank-1 SCFTs. Here authors also show that the assumption of the existence of certain rank-0 N=2 SCFTs leads to (eb) five additional consistent**

**rank-1 CB geometries. Subjects: High Energy Physics - Theory (hep-th) Cite as:arXiv: 1601.00011 [hep-th] (or arXiv: 1601.00011v2 [hep-th] for this version Geometric constraints on the space of N=2 SCFTs II: Construction of special Kähler geometries and RG flows Philip C. Argyres, Matteo Lotito, Yongchao Lü, Mario Martone (Submitted on 31 Dec 2015 (v1), last revised 6 Jan 2016 (this version, v2)) STOP Authors propose a model of Dark Energy in which the field currently dominating (e) the energy density of the universe is (=) an "axion field" linearly coupled to (e&eb) the Pontryagin density, tr (FΛF), (i.e., the exterior derivative of (e) the Chern-Simons form) of (e) a massive gauge field. Authors assume that the axion has (e) self-interactions corresponding to (e&eb) a non-trivial (exponential) potential. Authors argue that a non-vanishing magnetic helicity of (e) the gauge field triggers (eb) slow-rolling of the axion at field values far below (eb) the Planck scale. Proposal leads to (eb) a "Tracking Dark Energy Scenario" in which the contribution of the axion energy density to (e) the total energy density is (=) constant (and small) during (e&eb) the early radiation phase, until (e) a secular growth term proportional to (e) the Pontryagin density of (e) the gauge field becomes (=) dominant. The initially small contribution of (e) the axion field to (e) the total energy density is related to the observed small baryon-to-entropy ratio. Some speculations concerning the nature of the gauge field are offered.**

**Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Cite as:arXiv: 1601.00057 [hep-th] (or arXiv:1601.00057v2 [hep-th] for this version) Tracking Dark Energy from Axion-Gauge Field Couplings Stephon Alexander, Robert Brandenberger, Juerg Froehlich (Submitted on 1 Jan 2016 (v1), last revised 23 Jan 2016 (this version, v2)) STOP Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a**

**pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version) Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP In this work authors obtain bounds on (e&eb) the topological Abelian string-vortex and on (e&eb) the string-cigar, by using (e) a new measure of configurational complexity, known as (=) configurational entropy. In this way, the information-theoretical measure of (e) six-dimensional braneworld scenarios is capable to (e) probe situations where (e) the parameters responsible for (e) the brane thickness are arbitrary. The so-called configurational entropy (CE) selects (e&eb) the best value of the parameter in (eb) the model. This is accomplished by (e) minimizing the CE, namely, by (e) selecting the most appropriate parameters in the model that correspond to (e&eb) the most organized system, based upon (e) the Shannon information theory. This information-theoretical measure of complexity provides (eb) a complementary perspective to (e) situations where (e) strictly energy-based arguments are (=) inconclusive. Authors show that the higher the energy the higher the CE, what shows (eb) an important correlation between the energy of the localized field configuration and (e&eb) its**

**associated entropic measure. Subjects: High Energy Physics - Theory (hep-th) DOI:10.1016/j.physletb.2016.02.038 Cite as: arXiv: 1601.00076 [hep-th] (or arXiv: 1601.00076v2 [hep-th] for this version) Bounds on topological Abelian string-vortex and string-cigar from information-entropic measure R. A. C. Correa, D. M. Dantas, C. A. S. Almeida, Roldao da Rocha (Submitted on 1 Jan 2016 (v1), last revised 17 Feb 2016 (this version, v2)) STOP Sum of the Holevo quantity (that bounds the capacity of quantum channels to transmit (e&eb) classical information about an observable) and (e&eb) the quantum discord (a measure of the quantumness of (e) correlations of that observable) yields (eb) an observable-independent total given by (e) the quantum mutual information This split naturally delineates ,dilates and disseminates information about (e&eb) quantum systems accessible to (e) observers -- information that is redundantly transmitted by (e&eb) the environment -- while showing that it is (=) maximized for (e) the quasi-classical pointer observable. Other observables are accessible only via (e&eb) correlations with the pointer observable .Further authors prove an anti-symmetry property relating accessible information and (e&eb) discord. It shows that information becomes (=) objective -- accessible to (e) many observers -- only as quantum information is relegated to (e) correlations with (e&eb) the global environment, (eb) and, therefore, locally inaccessible. The resulting complementarity explains (eb) why, in a quantum Universe, we perceive objective classical reality while (e) flagrantly quantum superpositions are out of reach. Subjects: Quantum Physics (quant-ph) Journal reference: Scientific Reports 3, 1729 (2013) Report number: LAUR-13-26386 Cite as:arXiv: 1303.4659 [quant-ph] (or arXiv: 1303.4659v2 [quant-ph] for this version) Complementarity of quantum discord and classically accessible information Michael Zwolak, Wojciech H. Zurek STOP Niels Bohr proposed that the outcome of the measurement becomes (=) objective and real, and, hence, classical, when (e) its results can be communicated by (e) classical means. In this work authors revisit Bohr's postulate using (e) modern tools from (e) the quantum information theory. They find a**

**full confirmation of Bohr's idea: if a measurement device is in a nonclassical state, (eb) the measurement results cannot be communicated perfectly by classical means. In this case some part of information in the measurement apparatus is lost (e) in the process of communication: the amount of this lost information turns out to be (=) the quantum discord. The information loss occurs even when (e) the apparatus is not entangled with the system of interest. The tools presented in this work allow (eb) to generalize Bohr's postulate: authors show that for pure system-apparatus states (eb) quantum communication does not provide any advantage when (e) measurement results are communicated to (e) more than one recipient. They further demonstrate the superiority of quantum communication to (e&eb) two recipients on (e&eb) a mixed system-apparatus state and show that (eb) this effect is fundamentally different from (e) quantum state cloning. Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. Lett. 111, 040401 (2013) DOI:10.1103/PhysRevLett.111.040401 Report number: LA-UR-12-25351 Cite as: arXiv: 1304.2914 [quant-ph] (or arXiv: 1304.2914v2 [quant-ph] for this version) Quantum discord cannot be shared Alexander Streltsov, Wojciech H. Zurek STOP Wojciech H. Zurek, Don N. Page investigate thermodynamic equilibrium of (e) a self-gravitating perfect fluid in (eb) a spherically symmetric system containing (e) a black hole of mass M by means of (e) the Tolman-Oppenheimer-Volkoff (TOV) equation. At r >> 2M its solutions describe (eb) a black-body radiation atmosphere with (e&eb) the Hawking temperature T_BH~1/(8 \pi M) that is increasingly blueshifted as r approaches (eb) 2M. However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments:**

**Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by**

**(e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a firewall-like "Planck cocoon" with entropy close to (e) black hole entropy**

# Prescientiаl Dynamics And Contential Staticity Of Accentuation Attrition Models

## Module Four

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equailibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-445) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Some of the concepts to which model is applicable is as follows:**

**Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by (e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave signal. The signal sweeps upwards in frequency from 35 to 250 Hz with (e&eb) a peak gravitational-wave strain of (e) 1.0×10−21. It matches the waveform predicted by general relativity for (e) the inspiral and merger of a pair of black holes and the ringdown of the resulting (eb) single black hole. The signal was observed with a**

**matched-filter signal-to-noise ratio of (e) 24 and a false alarm rate estimated to be (=) less than 1 event per 203 000 years, equivalent to (=) a significance greater than 5.1σ. The source lies at (eb) a luminosity distance of 410+160−180 Mpc corresponding to (e&eb) a redshift z=0.09+0.03−0.04. In the source frame, the initial black hole masses are (=) 36+5−4M⊙ and 29+4−4M⊙, and the final black hole mass is (=) 62+4−4M⊙, with (e&eb) 3.0+0.5−0.5M⊙c2 radiated in (eb) gravitational waves. All uncertainties define (eb) 90% credible intervals. These observations demonstrate (eb) the existence of binary stellar-mass black hole systems. Observation of Gravitational Waves from a Binary Black Hole Merger B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 061102 – Published 11 February 2016 This is the first direct detection of gravitational waves and (e&eb) the first observation of a binary black hole merger. Received 21 January 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.061102 STOP This is the second in a series of three papers on (e&eb) systematic analysis of rank 1 Coulomb branch geometries of (e) four dimensional N=2 SCFTs. In the first paper authors developed a strategy for (e) classifying physical rank-1 CB geometries of N=2 SCFTs. They show how to carry out this strategy computationally to construct (eb) the Seiberg-Witten curves and (e&eb) one-forms for all the rank-1 SCFTs. Explicit expressions are given for all cases, with the exception of the Nf=4 SU (2) gauge theory and (e&eb) the En SCFTs which were previously constructed. Classification includes all known rank-1 theories plus (+) a new one with (e&eb) an abelian flavor group, plus (+) nine additional theories whose existence is more speculative. Four of those, reported in first paper, depend on (=) the assumption of new frozen rank-1 SCFTs. Here authors also show that the assumption of the existence of certain rank-0 N=2 SCFTs leads to (eb) five additional consistent rank-1 CB geometries. Subjects: High Energy Physics - Theory (hep-th) Cite as:arXiv: 1601.00011 [hep-th] (or arXiv: 1601.00011v2 [hep-th] for this version Geometric constraints on the space of N=2 SCFTs II: Construction of special Kähler geometries and RG flows Philip C. Argyres, Matteo Lotito, Yongchao Lü, Mario Martone (Submitted on 31 Dec 2015 (v1), last revised 6 Jan 2016 (this version, v2)) STOP Authors propose a model of Dark Energy in which the field currently dominating (e) the energy density of the universe is (=) an "axion field" linearly coupled to (e&eb) the Pontryagin density, tr (F∧F), (i.e., the exterior derivative of (e) the Chern-Simons form) of (e) a massive gauge field. Authors assume that the axion has (e) self-interactions corresponding to (e&eb) a non-trivial (exponential) potential. Authors argue that a non-vanishing magnetic helicity of (e) the gauge field triggers (eb) slow-rolling of the axion at field values far below (eb) the Planck scale. Proposal leads to (eb) a "Tracking Dark Energy Scenario" in which the contribution of the axion energy density to (e) the total energy density is (=) constant (and small) during (e&eb) the early radiation phase, until (e) a secular growth term proportional to (e) the Pontryagin density of (e) the gauge field becomes (=) dominant. The initially small contribution of (e) the axion field to (e) the total energy density is related to the observed small baryon-to-entropy ratio. Some speculations concerning the nature of the gauge field are offered.**
**Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Cite as:arXiv: 1601.00057 [hep-th] (or arXiv:1601.00057v2 [hep-th] for this version) Tracking Dark Energy from Axion-Gauge Field Couplings Stephon Alexander, Robert Brandenberger, Juerg Froehlich (Submitted on 1 Jan 2016 (v1), last revised 23 Jan 2016 (this version, v2)) STOP Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version) Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP In this work authors obtain bounds on (e&eb) the topological Abelian string-vortex and on (e&eb) the string-cigar, by using (e) a new**

**measure of configurational complexity, known as (=) configurational entropy. In this way, the information-theoretical measure of (e) six-dimensional braneworld scenarios is capable to (e) probe situations where (e) the parameters responsible for (e) the brane thickness are arbitrary. The so-called configurational entropy (CE) selects (e&eb) the best value of the parameter in (eb) the model. This is accomplished by (e) minimizing the CE, namely, by (e) selecting the most appropriate parameters in the model that correspond to (e&eb) the most organized system, based upon (e) the Shannon information theory. This information-theoretical measure of complexity provides (eb) a complementary perspective to (e) situations where (e) strictly energy-based arguments are (=) inconclusive. Authors show that the higher the energy the higher the CE, what shows (eb) an important correlation between the energy of the localized field configuration and (e&eb) its associated entropic measure. Subjects: High Energy Physics - Theory (hep-th) DOI:10.1016/j.physletb.2016.02.038 Cite as: arXiv: 1601.00076 [hep-th] (or arXiv: 1601.00076v2 [hep-th] for this version) Bounds on topological Abelian string-vortex and string-cigar from information-entropic measure R. A. C. Correa, D. M. Dantas, C. A. S. Almeida, Roldao da Rocha (Submitted on 1 Jan 2016 (v1), last revised 17 Feb 2016 (this version, v2)) STOP Sum of the Holevo quantity (that bounds the capacity of quantum channels to transmit (e&eb) classical information about an observable) and (e&eb) the quantum discord (a measure of the quantumness of (e) correlations of that observable) yields (eb) an observable-independent total given by (e) the quantum mutual information This split naturally delineates ,dilates and disseminates information about (e&eb) quantum systems accessible to (e) observers -- information that is redundantly transmitted by (e&eb) the environment -- while showing that it is (=) maximized for (e) the quasi-classical pointer observable. Other observables are accessible only via (e&eb) correlations with the pointer observable .Further authors prove an anti-symmetry property relating accessible information and (e&eb) discord. It shows that information becomes (=) objective -- accessible to (e) many observers -- only as quantum information is relegated to (e) correlations with (e&eb) the global environment, (eb) and, therefore, locally inaccessible. The resulting complementarity explains (eb) why, in a quantum Universe, we perceive objective classical reality while (e) flagrantly quantum superpositions are out of reach. Subjects: Quantum Physics (quant-ph) Journal reference: Scientific Reports 3, 1729 (2013) Report number: LAUR-13-26386 Cite as:arXiv: 1303.4659 [quant-ph] (or arXiv: 1303.4659v2 [quant-ph] for this version) Complementarity of quantum discord and classically accessible information Michael Zwolak, Wojciech H. Zurek STOP Niels Bohr proposed that the outcome of the measurement becomes (=) objective and real, and, hence, classical, when (e) its results can be communicated by (e) classical means. In this work authors revisit Bohr's postulate using (e) modern tools from (e) the quantum information theory. They find a full confirmation of Bohr's idea: if a measurement device is in a nonclassical state, (eb) the measurement results cannot be communicated perfectly by classical means. In this case some part of information in the measurement apparatus is lost (e) in the process of communication: the amount of this lost information turns out to be (=) the quantum discord. The information loss occurs even when (e) the apparatus is not entangled with the system of interest. The tools presented in this work allow (eb) to generalize Bohr's postulate: authors show that for pure system-apparatus states (eb) quantum communication does not provide any advantage when (e) measurement results are communicated to (e) more than one recipient. They further demonstrate the superiority of quantum communication to (e&eb) two recipients on (e&eb) a mixed system-apparatus state and show that (eb) this effect is fundamentally different from (e) quantum state cloning. Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. Lett. 111, 040401 (2013) DOI:10.1103/PhysRevLett.111.040401 Report number: LA-UR-12-25351 Cite as: arXiv: 1304.2914 [quant-ph] (or arXiv: 1304.2914v2 [quant-ph] for this version) Quantum discord cannot be shared Alexander Streltsov, Wojciech H. Zurek STOP Wojciech H. Zurek, Don N. Page investigate thermodynamic equilibrium of (e) a self-gravitating perfect fluid in (eb) a spherically symmetric system containing (e) a black hole of mass M by means of (e) the Tolman-Oppenheimer-Volkoff (TOV) equation. At r >> 2M its solutions describe (eb) a black-body radiation atmosphere with (e&eb) the Hawking temperature T_BH~1/(8 \pi M) that is increasingly blueshifted as r approaches (eb) 2M. However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments: Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by (e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a**

**firewall-like "Planck cocoon" with entropy close to (e) black hole entropy. Authors reproduce paper with an updated title and abstract. For a later study with similar results see G. 't Hooft, gr-qc/9706058 Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Physical Review D, Volume 29, Number 4, pp. 628 - 631 (1984) DOI:10.1103/PhysRevD.29.628 Cite as:arXiv:1511.07051 [gr-qc] (or arXiv:1511.07051v1 [gr-qc] for this version) Firewalls, black-hole thermodynamics, and singular solutions of the Tolman-Oppenheimer-Volkoff equation Wojciech H. Zurek, Don N. Page (Submitted on 22 Nov 2015) STOP Frederic Magniez, Ashwin Nayak consider the problem of testing the commutativity of (e) a black-box group specified by its k generators. The complexity (in terms of k) of (e) this problem was first considered by Pak, who gave (eb) a randomized algorithm involving O(k) group operations. They construct a quite optimal quantum algorithm for this problem whose complexity is in O~(k2/3)O~(k2/3). The algorithm uses and highlights the power of the quantization method of Szegedy. For the lower bound of Ω(k2/3)Ω(k2/3), authors give a reduction from (e) a special case of Element Distinctness to our problem. Along the way, they prove the optimality of the algorithm of (e) Pak for the randomized model. Algorithmica July 2007, Volume 48, Issue 3, pp 221-232 First online: 25 June 2007 Quantum Complexity of Testing Group Commutativity Frederic Magniez, Ashwin Nayak STOP Pawel Wocjan, Jon Yard analyze relationships between quantum computation and (e&eb) a family of generalizations of the Jones polynomial. Extending recent work by Aharonov et al., authors give efficient quantum circuits for (e) implementing the unitary Jones-Wenzl representations of the braid group. We use these to provide new quantum algorithms for (e) approximately evaluating a family of specializations of the HOMFLYPT two-variable polynomial of trace closures of braids. They also give (eb) algorithms for approximating the Jones polynomial of a general class of closures of (e) braids at roots of unity. Next authors provide (eb) a self-contained proof of a result of (e) Freedman et al. that any quantum computation can be replaced by (e&eb) an additive approximation of the Jones polynomial, evaluated at (eb) almost any primitive root of unity. Proof encodes two-qubit unitaries into (e&eb) the rectangular representation of (e) the eight-strand braid group. Authors then give QCMA-complete and (e&eb) PSPACE-complete problems which are based on braids. They conclude with direct proofs that evaluating the Jones polynomial of the plat closure at most primitive roots of unity is a #P-hard problem, while learning its most significant bit is PP-hard, circumventing (e) the usual route through the Tutte polynomial and graph coloring. Subjects: Quantum Physics (quant-ph); Geometric Topology (math.GT) Cite as: arXiv:quant-ph/0603069 (or arXiv:quant-ph/0603069v3 for this version) The Jones polynomial: quantum algorithms and applications in quantum complexity theory Pawel Wocjan, Jon Yard STOP Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan address the task of switching off the Hamiltonian of (e) a system by removing (e) all internal and system-environment couplings. Authors propose dynamical decoupling schemes, which use (e) only bounded-strength controls, for quantum many-body systems with local system Hamiltonians and local environmental couplings. To do so, they introduce (eb) the combinatorial concept of (e) balanced-cycle orthogonal arrays (BOAs) and show how to construct them from (e) classical error-correcting codes. The derived decoupling schemes may be useful as (=) a primitive for more complex schemes, e.g., for Hamiltonian simulation. For the case of n qubits and a 2-local Hamiltonian, the length of the resulting decoupling (e&eb) scheme scales as O (nlogn), improving over (e&eb) the previously best-known schemes that scaled quadratically with (e&eb) n. More generally, using balanced-cycle orthogonal arrays constructed from families of BCH codes, we show that bounded-strength decoupling for any ℓ-local Hamiltonian, where ℓ≥2, can be achieved using decoupling schemes of length at most O(nℓ−1logn). Subjects:Quantum Physics (quant-ph); Emerging Technologies (cs.ET) Journal reference: IEEE Trans. Inf. Theory, 62(5): 2881-2894, 2016 DOI: 10.1109/TIT.2016.2535183 Report number: MIT-CTP /4705 Cite as: arXiv: 1509.00408 [quant-ph] (or arXiv: 1509.00408v2 [quant-ph] for this version) Improved bounded-strength decoupling schemes for local Hamiltonians Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan STOP Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola propose dynamical control schemes for Hamiltonian simulation in many-body quantum systems that avoid instantaneous control operations and rely solely on (e) realistic bounded-strength control Hamiltonians. Each simulation protocol consists of (e) periodic repetitions of a basic control block, constructed as (=) a suitable modification of an "Eulerian decoupling cycle," that would otherwise implement (eb) a trivial (zero) target Hamiltonian. For an open quantum system coupled to (e) an uncontrollable environment, approach may be employed to (e) engineer an effective evolution that simulates (e&eb) a target Hamiltonian on the system, while suppressing (e) unwanted decoherence to the leading order. Authors present illustrative applications to (e) both closed- and open-system simulation settings, with (e&eb)**

**emphasis on simulation of non-local (two-body) Hamiltonians using (e) only local (one-body) controls. In particular, authors provide (eb) simulation schemes applicable to (e&eb) Heisenberg-coupled spin chains exposed to (e) general linear decoherence, and show how to simulate (e&eb) Kitaev's honeycomb lattice Hamiltonian starting from Ising-coupled qubits, as (=) potentially relevant to (e&eb) the dynamical generation of a topologically protected (e&eb) quantum memory. Additional implications for quantum information processing are discussed. Subjects: Quantum Physics (quant-ph) Journal reference: New J. Phys. 16, 045021 (2014) DOI: 10.1088/1367-2630/16/4/045021 Cite as:arXiv: 1310.4153 [quant-ph]**

**(or arXiv: 1310.4153v1 [quant-ph] for this version) Hamiltonian quantum simulation with bounded-strength controls Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola STOP Felix Fontein, Pawel Wocjan present a quantum algorithm for computing the period lattice of (e) infrastructures of fixed dimension. The algorithm applies to (e&eb) infrastructures that satisfy certain conditions. The latter are always fulfilled for (e) infrastructures obtained from (e) global fields, i.e., algebraic number fields and function fields with (e&eb) finite constant fields. The first main contributions is an exponentially better method for (e) sampling approximations of vectors of the dual lattice of (e) the period lattice than the methods outlined in the works of (e) Hallgren and Schmidt and Vollmer. This new method improves (eb) the success probability by (e) a factor of at least 2^ {n^2-1} where (e) n is the dimension. The second main contribution is a rigorous and complete proof that the running time of the algorithm is (=) polynomial in the logarithm of (e) the determinant of the period lattice and (e&eb) exponential in n. The third contribution is the determination of (e) an explicit lower bound on (e&eb) the success probability of our algorithm which greatly improves on (e&eb) the bounds given in the above works. The exponential scaling seems inevitable because (e) the best currently known methods for (e) carrying out fundamental arithmetic operations in infrastructures obtained from (e) algebraic number fields take exponential time. In contrast, the problem of computing (e&eb) the period lattice of infrastructures arising from function fields can be solved without (e) the exponential dependence on (e&eb) the dimension n since this problem reduces efficiently to the abelian hidden subgroup problem. This is also true for (e) other important computational problems in (eb) algebraic geometry. The running time of the best classical algorithms for (e) infrastructures arising from (e) global fields increases subexponentially with (e&eb) the determinant of the period lattice. Subjects: Quantum Physics (quant-ph) Cite as: arXiv: 1111.1348 [quant-ph]**

**(or arXiv: 1111.1348v2 [quant-ph] for this version) Quantum Algorithm for Computing the Period Lattice of an Infrastructure Felix Fontein, Pawel Wocjan STOP Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan advocate a new approach of (e) addressing hidden structure problems and (e&eb) finding efficient quantum algorithms. Authors introduce and investigate the Hidden Symmetry Subgroup Problem (HSSP), which is a generalization of the well-studied Hidden Subgroup Problem (HSP). Given a group acting on (e&eb) a set and an oracle whose level sets define (eb) a partition of the set; the task is to recover (e) the subgroup of symmetries of (e) this partition inside the group. The HSSP provides (eb) a unifying framework that, besides the HSP, encompasses (e) a wide range of algebraic oracle problems, including (e) quadratic hidden polynomial problems. While the HSSP can have (e) provably exponential quantum query complexity, authors obtain efficient quantum algorithms for various interesting cases. To achieve this, authors present a general method for (e) reducing the HSSP to the HSP, which works efficiently in several cases related to (e&eb) symmetries of polynomials. The HSSP therefore connects in a rather surprising way certain hidden polynomial problems with (e&eb) the HSP. Using this connection, authors obtain the first efficient quantum algorithm for (e) the hidden polynomial problem for (e) multivariate quadratic polynomials over fields of (e) constant characteristic. Authors also apply the new methods to (e&eb) polynomial function graph problems and (e&eb) present an efficient quantum procedure for (e) constant degree multivariate polynomials over (e&eb) any field. This result improves in several ways the currently known algorithms. Subjects: Quantum Physics (quant-ph) Journal reference: SIAM Journal on Computing 42:5 (2013) DOI: 10.1137/120864416 Cite as: arXiv: 1107.2189 [quant-ph]**

**(or arXiv: 1107.2189v2 [quant-ph] for this version) Hidden Symmetry Subgroup Problems Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan STOP Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan present an efficient general method for realizing a quantum walk operator corresponding to (e&eb) an arbitrary sparse classical random walk. Approach is based on Grover and Rudolph's method for preparing coherent versions of (e) efficiently integrable probability distributions. This method is intended for (e) use in quantum walk algorithms with polynomial speedups, whose complexity is usually measured in terms of (e&eb) how many times we have to apply a step of a quantum walk, compared to (e&eb) the number of necessary classical Markov chain steps. Authors consider a finer notion of complexity including (e) the**

**number of elementary gates it takes to implement each step of the quantum walk with (e&eb) some desired accuracy. The difference in complexity for (e) various implementation approaches is that method scales linearly in the sparsity parameter and (e&eb) poly-logarithmically with (e&eb) the inverse of the desired precision. Best previously known general methods either scale quadratically in (eb) the sparsity parameter, or polynomially in (eb) the inverse precision. Approach is especially relevant for (e) implementing quantum walks corresponding to (e&eb) classical random walks like those used in (eb) the classical algorithms for approximating permanents and sampling from binary contingency tables. In those algorithms, the sparsity parameter grows with (e&eb) the problem size, while (e) maintaining high precision is required. Subjects: Quantum Physics (quant-ph) Journal reference: Quantum Information and Computation, Vol.10, No.5&6, pp0420-0434 (2010) Cite as: arXiv: 0903.3465 [quant-ph] (or arXiv: 0903.3465v3 [quant-ph] for this version) Efficient Circuits for Quantum Walks Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan STOP Preparing the ground state of a system of interacting classical particles is (=) an NP-hard problem. Thus, there is in general no better algorithm to (e) solve this problem than exhaustively going through all N configurations of (e) the system to determine (eb) the one with lowest energy, requiring (e) a running time proportional to N. A quantum computer, if it could be built, could solve (e) this problem in time sqrt (N). Here, authors present a powerful extension of this result to (e) the case of interacting quantum particles, demonstrating (eb)that a quantum computer can prepare the ground state of (e) a quantum system as efficiently as it does for (e) classical systems. Subjects: Quantum Physics (quant-ph) Journal reference: Phys.Rev.Lett.102:130503, 2009 DOI:10.1103/PhysRevLett.102.130503 Cite as: arXiv: 0809.2705 [quant-ph] (or arXiv: 0809.2705v1 [quant-ph] for this version) Preparing ground states of quantum many-body systems on a quantum computer David Poulin, Pawel Wocjan STOP Daniel Nagaj, Pawel Wocjan construct a simple translationally invariant, nearest-neighbor Hamiltonian on (e&eb) a chain of 10-dimensional qudits that makes it possible to (e) realize universal quantum computing without (e) any external control during (e&eb) the computational process. Authors only require the ability to prepare an initial computational basis state which encodes (e) both the quantum circuit and its input. The computational process is (=) then carried out by (e) the autonomous Hamiltonian time evolution. After a time polynomially long in the size of (e) the quantum circuit has passed, the result of (e) the computation is obtained with (e&eb) high probability by (e) measuring a few qudits in the computational basis. This result also implies (eb) that there cannot (e) exist efficient classical simulation methods for (e) generic translationally invariant nearest-neighbor Hamiltonians on (e&eb) qudit chains, unless (e) quantum computers can be efficiently simulated by (e&eb) classical computers (or, put in complexity theoretic terms, unless BPP=BQP). Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. A 78, 032311 (2008) DOI: 10.1103/PhysRevA.78.032311 Cite as: arXiv: 0802.0886 [quant-ph] (or arXiv: 0802.0886v4 [quant-ph] for this version) Hamiltonian Quantum Cellular Automata in 1D Daniel Nagaj, Pawel Wocjan**

## 19. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_4}{dt} = -b_4 Q_4 + g_4(g_4)Q_4$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_4(0) = 0, \quad g_4(P_4) > 0, \; P_4 \; 0, \qquad \lim_{P_4 \to \infty} g_4(P_4) = (K_4)_4 \; \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_4 + (\rho_a)_4 + D_4(a_4, P_4, Q_4)\rho_4 + 0, \qquad 0 < a_4 \; \infty,$$

$$\rho_4(t, 0) = \int_0^\infty n_4(a_4, P_4, Q_4)\rho_4(t_4, a_4)da_4$$

Where $D_4$ is the mortality function and $n_4$ is the fertility function.

Put $_4$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_4}{dt}\right]_4 = x_4\left[\frac{dQ_2}{dt}\right]_4 - x_2\left[\frac{dQ_4}{dt}\right]_4$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_4 - x_4\left[\frac{dQ_2}{dt}\right]_4$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_4 - x_6\left[\frac{dQ_3}{dt}\right]_4$$

Where $x_4, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_4$ and $g_4(P_4)Q_4$

Now when you add the term $(b_4Q_4)$ with $g_4(P_4)(Q_4)$ we are infect see is that $(b_4Q_4)$ is credited incrementally part of $g_4(P_4)(Q_4)$ or instant nearly, when subtracted from $\frac{dQ_4}{dt}$, it means that $g_4(P_4)(Q_4)$ is slowly eaten by $\frac{dQ_4}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_4 = \mu_4(a_4) + (f_4)_4(P_4)_4 + (f_2)_4(Q)_4 \quad (2)$$

$$n_4 = \beta_4(a_4)h_4(Q_4), \quad (3)$$

Where $\mu_4$the age- dependent debit rate whit is $\mu_4(a_4)$ $\forall a > 0$; $(f_4)_4$ is the density- dependent death rate with $f_4(0) = 0$, $(f_4^4)_4(P_4)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_4)_4(P_4) = (k_4)_4$ ; $f_2$ is the functional response of creditors to prey with $[f_2]_4(0) = 0, (f_2^4)_4(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_4(Q_4) = [k_2]$ ; $\beta$ is the age dependent birth rate with $\beta(a)$ $0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 4, h^4(Q)$ $0, \forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_4 + (\rho_a)_4 = -\mu_4(a_4) + (f_4)_4P_4) + \big(f_2)_4(Q_4)\big)(\rho_4)_4 \quad (4)$$

$$\rho_4(t_4, 0) = h_4(Q_4)\int_0^\infty \beta_4(a_4)\rho_4(t, a_4)da, \quad (5)$$

$$\frac{dQ_4}{dt}$$

$$= -b_4Q_4 + g_4(P_4)Q_4 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_4$ and $Q_4$ of (4-6) satisfy

$$0 \leq \rho_4(t, a_4) \leq \rho_4^*(t, a_4),$$

$$0 \leq Q_4(t) \leq Q_4^*(t),$$

Where $\rho_4^*(t, a_4)$ and $Q_4^*(t)$, are solutions of

$$(\rho_t^*)_4 + (\rho_a^*)_4 = -(\mu_4(a_4) + (f_4)_4(P_4^*) + (f_2)_4 Q_4^*))\, \rho_4^*, \quad (7)$$

$$\rho_4^*(t, 0) = \int_0^\infty \beta_4(a_4)\, \rho_4^*(t, a_4) da_4 \quad (8)$$

$$\frac{dQ_4^*}{dt} = -b_4 Q_4^* + g_4(P_4^*)Q_4^*, \quad (9)$$

With $P_4^*(t) = \int_0^\infty \rho_4^*(t, a_4) da_4$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_4^*(t)}{dt} = P_4^*(t)\big(r_4 - (f_4)(P_4^*) - (f_2)_4(Q^*)\big), \quad (10)$$

$$\frac{dQ_4^*(t)}{dt} = Q_4^*(t)(-b_4 + g_4(P_4^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_4(a) e^{-\int_0^a \mu_4(\tau)dt} e^{-ra}\, da_4 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_4^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^{\infty} \beta_4(a_4) e^{-\int_0^a \mu_4(t)dt}\, da_4 > 1, \qquad (H1)$$

Which guarantees $r_4 > 0$.

*Remark.* If $\beta_4(a) = [\beta_0]_4 a_4 e^{-\alpha_4 a_4}, \mu_4(a_4) \equiv \mu_4$ then the solution of (12) is

$$r_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_4]_4(P_4) = 0, [f_2](Q_4) = \lambda_4 Q_4\ h_4(Q_4) = 1, \qquad and\ \ g_4(P_4) = c_4 P_4$$

$$[\rho_4]_4 + [\rho_a]_4 = -(\mu_4 + \lambda_4 Q_4)\rho_4$$

$$\rho_4(t,0) = \int_0^{\infty} [\beta_0]_4 a_4 e^{-\alpha_4 a_4} \rho_4(t, a_4)\, da_4$$

$$\frac{dQ_4}{dt} = -b_4 Q_4 + c_4 P_4 Q_4$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_4}{dt} = \omega_4 P_4 - \lambda_4 P_4 Q_4$$

$$\frac{dQ_4}{dt} = b_4 Q_4 - c_4 P_4 Q_4$$

Where $\omega_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$

## 20. THE EQUILIBRIUM $(\rho_0)_4(a_4), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_4(a_4), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_4 + (\rho_a)_4 = -\big(\mu_4(a_4) + (f_4)_4(P_4)\big)\rho,$$

$$\rho_4(t,0) = \int_0^\infty \beta_4(a_4)\rho_4(t,a_4)da_4 \qquad (13)$$

The solution of (13) has the form of

$$\rho_4(t,a_4) = Ce_4^{r(t-a_4)} - \int_0^a (\tau)d\tau - \int_0^4 (f_4)_4\big(P_4(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_4(t) = C_4 e^{r_4 t - \int_0^t (f_4)_4(P_4(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_4^{-4}(r)$ if $k_4 > r$. This is equivalent to

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{(k_4)_4 a_4} da_4 < 1.$$

In fact, if we define

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau)d\tau} e^{-s_4 a_4}\, da_4$$

$R(s)$ is a decreasing function. $R_4(r_4) = 1$ And $R_4(k_4)_4 < 1$ imply $r_4 < (k_4)_4$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_4(a_4), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{k_4 a_4} da_4$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_4 \equiv \int_0^\infty (\rho_0)_{44}(a_4)da_4$ *satisfies*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{-(f_4)_4(P_0)a_4} da$$

$$= 1 \qquad (11)$$

Next, we study the stability of $(\rho_0)_4(a_4), 0)$

Let $u_4(t, a_4) = \rho(t, a_4) - \rho_0(a_4)$ Then

$$N_4(t) \equiv \int_0^\infty u_4(t, a_4)da_4 = \int_0^\infty \rho_4(t, a_4)da_4 - \int_0^\infty (\rho_0)_4(a_4)da = P_4(t) - P_0)_4$$

And

$$(u_4)_4 + (u_a)_4 = (\rho_t)_4 + (\rho_a)_4 - (\rho_0^4)_4(a_4)$$

$$= -\mu_4(a_4) + (f_4)_4(P_4) + (f_2)_4(Q_4))\,(u_4 + (\rho_0)_4) + \mu_4(a_4)(\rho_0)_4 + (f_4)_4(P_0)_4(\rho_0)_4$$

$$= -\mu_4(a_4)u_4 - (f_4)_4(P_4) + (f_2)_4(Q_4))\,u_4 - (\rho_0)_4\,(f_4)_4(P_4) - (f_4)_4(P_0)_4$$

$$+ (f_2)_4(Q_4)), \qquad (15)$$

Where $(\rho_0^4)_4(a_4)$ denotes $d(\rho_0)_4(a_4)/da_4$

Expansion $(f_4)_4(P_4)$ and $(f_2)_4(Q_4)$ at $(P_0)_4$ and 0 yield respectively

$$(f_4)_4(P_4) = (f_4)_4((P_0)_4 + N_4)$$

$$= (f_4)_4(P_0)_4 + (f_4^4)_4(P_0)_4 N_4 + o_4(N_4), \quad (16)$$

$$(f_2)_4(Q_4)$$

$$= (f_2^4)_4(0)Q + o_4(Q_4) \quad (17)$$

Write $f_4(P_0) = f_{40}, f_4'(P_0) = f_{40}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_4)_4 + (u_a)_4 \approx -(\mu_4(a_4) + (f_{40})_4)u - (f_{40}^4)_4(\rho_0)_4(N_4) - (f_{20}^4)_4(\rho_0)_4(Q_4)$$

$$u_4(t,0) = \rho_4(t,0)(\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a) + (\rho_0)_4(a_4)da_4 - (\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a_4)da_4 + (h_4(Q_4) - 4)(\rho_0)_4(0) \quad (18)$$

Expanding

$$h_4(Q_4) = 1 + (h^4)_4(0)Q_4 + 0(Q_4) \equiv 1 + h_0'Q_4 + 0Q_4$$

And substituting it into (48), we arrive at

$$u_4(t,0) \approx \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + (h_{a_4}')(\rho_0)_4 Q_4$$

The linearization is the following :

$$[u_4] + [u_4] = -(\mu_4(a_4) + [t_{40}]_4 u_4 - [\alpha_4]_4(a_4)N_4 - [\alpha_2]_4(a_4)Q_4$$

$$u_4(t,0) = \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + [\xi_4]Q_4 \quad (19)$$

$$\frac{dQ_4}{dt} = -b_4 Q_4$$

where

$$[\alpha_4]_4(a_4) = [t_4^4 a]_4[p, o]_4(a_4) d2(a_4) = [f_{20}^4]_4(\rho\, o\,)_4(a_4)\xi = [\mathrm{h\ o}^4]_4[\rho \mathrm{o}]_4(\mathrm{o}) \text{ and } \mathrm{b}_4 = \mathrm{b}_4 - g_4(P_4)$$

If $g_4(P_0) > b$, clearly, $(\rho_0)_4(a_4), 0)$ is unstable and the predator population is persistent. Assume $g_4(P_0)_4 < b$.

Then, $Q_4(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_4 + u_a = -(\mu_4(a_4) + (f_{40})_4)u_4 - (\alpha_4)_4(a_4)N_4$$

$$= -(\mu_4(a_4) + f_{40})u_4 - \alpha_4(a_4)\int_0^\infty u_4(t, a_4)da_4$$

$$u_4(t, 0) = \int_0^\infty \beta_4\,(a_4)u_4(t, a_4)da_4$$

Using the method developed by Cushing [2, 3], we suppose

$$u_4(t, a_4) = \tilde{\mathrm{u}}_4(a_4)^{e^{\gamma(t-a_4)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_4}\, da_4$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_4)$ gives,

$$\tilde{\mathrm{u}}_4(a_4) = e^{-\int_0^a \mu_4(\tau_4)d\tau - (f_{40})_4 a}\left(\tilde{\mathrm{u}}_4(0) - w\int_0^a e^{\int_0^\tau \mu_4(s_4)ds + (f_{40})}\tau_4\,(\alpha_4)_4(\tau_4)e^{\gamma_4\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1-\int_0^{\infty}\beta_4(a_4)e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\right)\tilde{u}_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a(\alpha_4)_4(\tau_4)e^{-\gamma_4\tau_4}e^{(f_{40})_4\tau_4+\int_0^\tau \mu_4(s_4)ds}d\tau_4\,da_4$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}da\right)(\tilde{u}_4(0))_4$$

$$+\left(-4-\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a e^{(f_{40})_4\tau_4+\int_0^\tau \mu_4(s_4)ds+y\tau}\alpha_4(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(t_4)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_4\,(a_4)Y_4(\gamma_4,a_4)da_4\right)u_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)\int_0^a(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4=0,$$

(24)

$$\left(\int_0^{\infty}Y_4(\gamma_4,a_4)da_4\right)\tilde{u}_4(0)$$

$$+\left(-1-\int_0^{\infty} Y_4(\gamma_4,a_4)\int_0^{a}(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4\right)w=0.$$

Equqtion (20) has a solution of the form of $u_4(t,a_4)=u_4(a_4)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_4(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_4(\gamma_4,a_4)=\int_0^a(\rho_0)_4(\tau_4)Y^{-4}(\gamma_4,a_4)d\tau_4$. We obtain the characteristic equation

$$f_{40}^{l}{}_4\left(\int_0^{\infty} Y_4(\gamma_4,a_4)da_4\right)\left(\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$+\left(1-\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)da\right)\left(1+(f_{40}^{l})_4\int_0^{\infty}Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_4(a_4)=(\rho_0)_4(0)e^{-\int_0^a\mu_4(\tau_4)d\tau_4-(f_{40})_4a_4}$$

$$I_4(\gamma_4,a_4)=\int_0^a(\rho_0)_4(\tau_4)e^{\gamma_4\tau_4+(f_{40})_4\tau+\int_0^{\tau}\mu_4(s_4)ds_4}d\tau$$

$$=\frac{(\rho_0)_4(0)}{\gamma_4}(e^{\gamma_4a_4}-1)$$

Let

$$E_4(a_4)=e^{-(f_{40})_4a_4-\int_0^a\mu_4(\tau_4)d\tau_4}$$

Then $[Y_4]_4[\gamma_4]_4, a_4) = E_4(a_4)e^{-\gamma_4 a_4}$. by a simple calculation,

$$\int_0^\infty \beta_4(a_4)Y_4(\gamma_4, a_4)I(\gamma_4, a_4)da = \frac{[\rho_0]_4(0)}{\gamma}\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}\ da\right);$$

$$\int_0^\infty Y_4(\gamma_4, a_4)I(\gamma_4, a_4)da = \frac{[\rho_0]_4(0)}{\gamma_4}\left(\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}\ da_4\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_4(0)}{\gamma_4}[f_{40}^4]\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}da_4\ (\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4)$$

$$+\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right).$$

$$\left(1 + \frac{(\rho_0)_4(0)}{\gamma_4}(f_{40}^{'})_4\left(\int_0^\infty E_4(\,a_4)da_4 - \int_0^\infty E_4(\,a_4)e^{-\gamma_4 a_4}da_4\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right)(1 + \frac{(f_{40}^{'})_4}{\gamma_4}(\rho_0)_4(0)\int_0^\infty E_4(\,a_4)da_4) = 0.$$

We need ũ(a) $\not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1 - \int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)da_4 = 1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4 \neq 0,$$

Since $\int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)\ \int_0^a \alpha_4\,(\tau_4)Y^{-4}(\gamma_4, \tau_4)d\tau_4 da_4 > 0$. Therefore, the characteristic solution is

$$\gamma_4 = -(f_{40}^{'})_4(\rho_0)_4 \int_0^{\infty} E_4(a_4)da_4.$$

Since $f_4'(P_0)_4 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_4), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_4), 0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_4 + (\mu_4(a_4) + [f_4]_4(P_4^*) + [f_2]_4(Q_4^*))P_4^* = 0, \quad (26)$$

$$p_4^*(0) = h_4(Q_4^*)\int_0^{\infty} \beta_4(a_4)p_4^*(a_J)p_4^*(a_4)da_4 \quad (27)$$

$$g_4(P_4^*) - b_4 = 0, \quad (28)$$

Where $P_4^* = \int_0^{\infty} p_4^*(a_4)\, da_4$

Equation (28) yields $P^* = g^{-4}(b)$,

From (26),

$p_4^*(a) = p_4^*(0)e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}$

Substituting into (27), we need to solve

$$\left(1 - h_4(Q_4^*)\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}da_4\right)p_4^*(0) = 0$$

For $Q_4^*$

Fix $p_4^*(0)$ and $P_4^*$ Let

$$H_4(z_4) = 1 - h_4(z_4)\int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}da_{,4}$$

And assume that

$$[k_3]_4\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(g_4^{-4}(b_4))a_4 - [k_2]_4 a_4}\, da_4 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_4(z_4)] > 0.$$

Clearly,

$$\frac{dH_4}{dz_4} = h_4(z)(\,[f_2']_4(z_4)\,)\int_0^\infty a_4\beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - f_4(P_4^*)a_4 - f_2(z_4)a_4}da_{\square}$$

$-h'(z_4)\int_0^\infty \beta_4(a_4)\, e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}\quad da_4$

the assumptions on $\left[f_2^4\right]_4 and\ [h^4]^4 dH_4/dz > 0$. Hence, if

$$H_4(0) = 1 - \int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4} da_4 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_4(0) < 0$ follows if $g_4^-(b_4) < [f_4^{-4}]_4(r_4)$, where $r_4$ is the solution of (12).

Indeed, since

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 e^{-y\frac{a}{4}} da_4}$$

Is a decreasing function of $s_4$ when $P_4^* = [g^{-4}]_4(b_4) < [f_4^{-4}]_4(r_4)$, that is $[f_4]_4(P_4^*) < r$, $R\left(f_4(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_4(g^{-4}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4\left(g_4^{-4}(b)\right)a_4} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_4^*(a_4), Q_4^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_4^*(a_4), Q_4^*)$ is more complicated. We again consider perturbations $m_4(t_4 a_4) = p_4(t_4, a_4) - p_4^*(a_4)$ and $W_4(t) = Q_4(t) - Q^*$.Then

$$M_4(t) \equiv \int_0^\infty m_4(t, a_4) da_4 = P_4(t) - P_4^*,$$

$$[m_4]_4 + [m_a]_4 = [p_4]_4 + [p_a]_4 - [p^*]_4(a_4) \approx -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4$$

$$-[f_{4*}^4]_4 p_4^*(a_4)M - [f_{2*}^4]_4 p_4^*(\mathrm{a})W_4,$$

$$m_4(t,0)p_4(t,0) - p_4^*(0) \approx h_* \int_0^\infty \beta_4(a_4)m_4(t,a_4)\,da + \frac{[h_*^4]_4}{[h_*]_4}\, p_4^*(0)W_4$$

$$\frac{dW_4}{dt} = \frac{dQ_4}{dt} \approx [g_*^4]_4 Q_4^* M_{,4}$$

Where $[f_{4*}]_4$ denoting $[f_4(P^*)]_4$, $[f_{4*}]_4$ denoting $[f_2(P^*)]_4$, $[f_{4*}^{'}]_4$ denoting $[f^{'}{}_4(P^*)]_4 [f_{2*}^{'}]_4$ denoting $[f_2]_4$ $(Q_4^*,$ $h_{*}$ denoting $h_*^{'}$ denoting $h_4^{'}$ $(Q_4^*)$, and $[g_*^{'}]_4$

Denoting $[g^{'}]_4(P_4^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_4]_4 + [m_a]_4 = -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4 - [f^{'}{}_{4*}]_4 [p_4^*](a_4)M_4 - [f_{2*}^{'}]_4 p_4^*(a_4)W_4$$

$$m_4(t,0) = [h_*]_4 \int_0^\infty \beta_4(a_4)m_4(t,a_4)da_4 + \frac{[h_*^4]_4}{[h^*]_4}[p^*]_4(0)\, W_{,4} \qquad (30)$$

$$\frac{dW_4}{dt} = [g_*^4]_4 [Q_4^*]M_4$$

This is equation is important in the following cases: problem of using observed correlations to infer causal relations is relevant to (e) a wide variety of scientific disciplines. Yet given correlations between just two classical variables, (eb) it is impossible to determine whether they arose from (e) a causal influence of one on the other or a common cause influencing both, unless one can implement (eb) a randomized intervention. **Katja Ried, Megan Agnew,**

**Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. Resch** here consider the problem of causal inference for (e) quantum variables. Authors introduce causal tomography, which unifies and generalizes (e&eb) conventional quantum tomography schemes to provide (eb) a complete solution to the causal inference problem using (e) a quantum analogue of a randomized trial. Authors furthermore show that, in contrast to the classical case, observed quantum correlations alone can sometimes provide (eb) a solution. They implement a quantum-optical experiment that allows (eb) to control (e&eb) the causal relation between two optical modes, and two measurement schemes -- one with (e&eb) and one without (e) randomization -- that extract this relation from (e) the observed correlations. Results show that entanglement and (e&eb) coherence, known to be central to (e) quantum information processing, also provide (eb) a quantum advantage for (e) causal inference. Subjects: Quantum Physics (quant-ph); Learning (cs.LG); General Relativity and Quantum Cosmology (gr-qc); Machine Learning (stat.ML) Journal reference: Nat Phys 11, 414-420 (2015) DOI: 10.1038/nphys3266 Cite as: arXiv: 1406.5036 [quant-ph] (or arXiv: 1406.5036v1 [quant-ph] for this version) **Inferring causal structure: a quantum advantage Katja Ried, Megan Agnew, Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. ReschSTOP D. T. Son, A. O. Starinets** review recent progress in applying the AdS/CFT correspondence to (e&eb) finite-temperature field theory. In particular, they show how the hydrodynamic behavior of field theory is reflected in (e&eb) the low-momentum limit of (e) correlation functions computed through (e&eb) a real-time AdS/CFT prescription, which they formulate. Authors also show how the hydrodynamic modes in field theory correspond to (e&eb) the low-lying quasinormal modes of (e) the AdS black p-brane metric. They provide a proof of the universality of (e) the viscosity/entropy ratio within (eb) a class of theories with (e&eb) gravity duals and formulate (eb) a viscosity bound conjecture. Possible implications for real systems are mentioned. Subjects: High Energy Physics - Theory (hep-th) Journal reference: Ann.Rev.Nucl.Part.Sci.57:95-118,2007 DOI: 10.1146/annurev.nucl.57.090506.123120Report number: INT PUB 07-02Cite as: arXiv:0704.0240 [hep-th]

(or arXiv:0704.0240v2 [hep-th] for this version) An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society **Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP** Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv:

1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version **Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP Viscosity, Black Holes, and Quantum Field Theory D. T. Son, A. O. Starinets STOP Authors** study information retrieval from evaporating black holes, assuming that the internal dynamics of a black hole is unitary and rapidly mixing, and assuming that the retriever has unlimited control over the emitted Hawking radiation. If the evaporation of the black hole has already proceeded past the ``half-way" point, where (e) half of the initial entropy has been radiated away, then (eb) additional quantum information deposited in (eb) the black hole is revealed in (eb) the Hawking radiation very rapidly. Information deposited prior to the half-way point remains concealed until (e) the half-way point, and then emerges (eb) quickly. These conclusions hold because (e) typical local quantum circuits are efficient encoders for (e) quantum error-correcting codes that nearly achieve (eb) the capacity of the quantum erasure channel. Estimate of a black hole's information retention time, based on (e) speculative dynamical assumptions, is just barely compatible with (e&eb) the black hole complementarity hypothesis. **Black holes as mirrors: quantum information in random subsystems Patrick Hayden1 and John Preskill2 Published 26 September 2007 Journal of High Energy Physics, Volume 2007, JHEP09 (2007) STOP** By analyzing the key properties of black holes from the point of (e) view of quantum information, authors derive a model-independent picture of (e) black hole quantum computing. It has been noticed that this picture exhibits (eb) striking similarities with (e&eb) quantum critical condensates, allowing (eb) the use of a common language to describe (eb) quantum computing in both systems. Authors analyze such quantum computing by allowing coupling to (e&eb) external modes, under the condition that the external influence must be (=) soft-enough in order not to (e) offset the basic properties of the system. Authors derive model-independent bounds on some crucial time-scales, such as (=) the times of gate operation, decoherence, maximal entanglement and total scrambling. They show that for black hole type quantum computers all these time-scales are of (e) the order of the

black hole half-life time. Furthermore, authors construct explicitly a set of Hamiltonians that generates (eb) a universal set of quantum gates for (e) the black hole type computer. They find that the gates work at maximal energy efficiency. Furthermore, authors establish a fundamental bound on the complexity of quantum circuits encoded on (eb&e) these systems, and characterize the unitary operations that are (=) implementable. It becomes apparent that the computational power is very limited due to (e) the fact that the black hole life-time is of the same order of (e) the gate operation time. As a consequence, it is impossible to retrieve its information, within (eb) the life-time of a black hole, by externally coupling to (e&eb) the black hole qubits. However, authors show that, in principle, coupling to (e&eb) some of the internal degrees of freedom allows (eb) acquiring knowledge about the micro-state. Still, due to the trivial complexity of operations that can be performed, (eb) there is no time advantage over the collection of Hawking radiation and subsequent decoding. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Report number: LMU-ASC 03/16, MPP-2016-1 Cite as: arXiv:1605.01407 [hep-th] (or arXiv:1605.01407v1 [hep-th] for this version) **Universality of Black Hole Quantum Computing Gia Dvali, Cesar Gomez, Dieter Lust, Yasser Omar, and Benedikt Richter (Submitted on 4 May 2016) STOP** In this addendum to [arXiv: 1402.5674] two points are discussed. In the first additional evidence is provided (eb) for a dual connection between the geometric length of an Einstein-Rosen bridge and (e&eb) the computational complexity of the quantum state of the dual CFT's. The relation between growth of complexity and (e&eb) Page's “Extreme Cosmic Censorship” principle is also remarked on. The second point involves a gedanken experiment in which Alice measures (eb) a complete set of commuting observables at her end of (e) an Einstein-Rosen bridge is discussed. An apparent paradox is resolved by appealing to (e) the properties of GHZ tripartite entanglement. **Addendum to computational complexity and black hole horizons Leonard Susskind* Version of Record online: 18 JAN 2016 DOI: 10.1002/prop.201500093 STOP** Analyzing some well established facts, authors give a model-independent

parameterization of black hole quantum computing in terms of (e&*eb) a set of macro and micro quantities and their relations. These include the relations between the extraordinarily-small energy gap of (e) black hole qubits and important time-scales of (e) information-processing, such as, scrambling time and Page's time. Authors then show, confirming and extending previous results, that other systems of nature with (e&eb) identical quantum informatics features are (=) attractive Bose-Einstein systems at (eb) the critical point of quantum phase transition. Here authors establish a complete isomorphy between (e&eb) the quantum computational properties of these two systems. In particular, they show that the quantum hair of a critical condensate is (=) strikingly similar to (=) the quantum hair of a black hole. Irrespectively whether one takes the similarity between the two systems as a remarkable coincidence or as (=) a sign of a deeper underlying connection, the following is evident. Black holes are not unique in their way of (e) quantum information processing and they can manufacture black hole based quantum computers in labs by (e) taking advantage of quantum criticality. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Atomic Physics (physics. Atom-ph); Quantum Physics (quant-ph) Cite as: arXiv: 1601.01329 [hep-th] (or arXiv: 1601.01329v1 [hep-th] for this version) **Black Hole Based Quantum Computing in Labs and in the Sky Gia Dvali, Mischa Panchenko (Submitted on 6 Jan 2016) STOP Gia Dvali, Alexander Gußmann** show that the existence of black holes with (e&eb) classical skyrmion hair invalidates (e) standard proofs that global charges, such as (=) the baryon number, cannot be conserved by (e) a black hole. By carefully analyzing the standard arguments based on (e) a Gedankenexperiment in which a black hole is seemingly-unable to (e) return the baryon number that it swallowed, authors identify inconsistencies in this reasoning, which does not (e) take into the account neither the existence of skyrmion black holes nor the baryon/skyrmion correspondence. Authors then perform a refined Gedankenexperiment by incorporating (e) the new knowledge and show that (eb) no contradiction with conservation of baryon number takes place at (eb) any stage of black hole evolution. Analysis also indicates (eb) no conflict

between semi-classical black holes and (e&eb) the existence of baryonic gauge interaction arbitrarily-weaker than gravity. Next, authors study classical cross sections of a minimally-coupled massless probe scalar field scattered by (e) _ a skyrmion black hole. Authors investigate how the skyrmion hair manifests itself by (e) comparing this cross section with (e&eb) the analogous cross section caused by (e) a Schwarzschild black hole which has (e) the same ADM mass as (=) the skyrmion black hole. Here authors find an order-one difference in the positions of (e) the characteristic peaks in (eb) the cross sections. The peaks are shifted to smaller scattering angles when (e) the skyrmion hair is present. This comes from the fact that the skyrmion hair changes (e&eb) the near horizon geometry of the black hole when compared to (e&eb) a Schwarzschild black hole with same ADM mass. Authors keep the study of this second aspect general so that (e) the qualitative results obtained can also be applied to(e&eb) black holes with classical hair of different kind. Subjects: High Energy Physics - Theory (hep-th); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU-ASC 18/16 Cite as: arXiv:1605.00543 [hep-th] (or arXiv:1605.00543v2 [hep-th] for this version) A preliminary discussion is given of (e) the prospects that gravitational-wave observations of (e) binary inspiral of black holes could reveal or constrain (e&eb) quantum modifications to (e) black hole dynamics, such as are required to (e) preserve postulates of quantum mechanics. Different proposals for such modifications are characterized by (e) different scales, and the size of these scales relative to (e&eb) those probed by observation of (e) inspiral signals is important in determining (eb) the feasibility of finding experimental signatures. Certain scenarios with strong quantum modifications in (eb) a region extending well outside the horizon are expected to modify (e&eb) classical evolution, and distort (e&eb) the near-peak gravitational wave signal, suggesting (eb) a search for departures from (e) waveforms predicted by general relativity. The near agreement of the GW150914 signals with (e&eb) such waveforms is discussed, and indicates (eb) constraints on some such scenarios. Important strategies for more precise future tests are 1) to develop more precise

predictions from (e) scenarios proposing quantum modifications, and 2) searching for observed deviations from (e) numerical relativity predictions via (e&eb) analysis of gravity wave data, particularly focussing on (e&eb) the signal region corresponding to plunge and merger. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1602.03622 [gr-qc] (or arXiv: 1602.03622v2 [gr-qc] for this version) **Gravitational wave tests of quantum modifications to black hole structure -- with post-GW150914 update Steven B. Giddings STOP** Direct LHC bounds on (e&eb) colored SUSY particles now corner naturalness more than the measured value of the Higgs mass does. Bounds on the gluino are of (e) particular importance, since it radiatively "sucks" up (e) the stop and Higgs-up soft masses. As a result, even models that easily accommodate a 125 GeV Higgs are almost as tuned as (=) the simplest version of SUSY, the MSSM: at best at the percent level. In this paper, authors further examine how current LHC results constrain (e) naturalness in three classes of models that may relax LHC bounds on (e&eb) sparticles: split families (e&eb), baryonic RPV (e&eb), and Dirac gauginos. In models of split families and bRPV, the bounds on the gluino are (=) only slightly reduced, resulting in (eb) a few percent tuning. In particular, having a natural spectrum in bRPV models typically implies (eb) that tops, Ws, and Zs are easily produced (eb) in the cascade decays of (e) squarks and gluinos. The resulting leptons and missing energy push (eb) the gluino mass limit above 1 TeV. Even when the gluino has a Dirac mass and does not contribute to the stop mass at one loop, tuning reappears in calculable models because there is no symmetry imposing the supersoft limit. Authors conclude that, even if sparticles are found at LHC-14, naturalness will not emerge triumphant. Subjects: High Energy Physics - Phenomenology (hep-ph)DOI:10.1007/JHEP03(2014)022 Cite as:rXiv:1309.3568 [hep-ph] (or arXiv:1309.3568v2 [hep-ph] for this version) **The Last Vestiges of Naturalness Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang, Ken Van Tilburg, Giovanni Villadoro STOP** The lack of evidence for new physics beyond (e) the standard model at the LHC points to (e) a paucity of new particles near the weak scale. This suggests that the weak scale is tuned and that supersymmetry, if present at

all, is realized (eb) at higher energies. The measured Higgs mass constrains (e) the scalar sparticles to be below 10^5 TeV, while (e) gauge coupling unification favors (eb) Higgsinos below 100 TeV. Nevertheless, in many models gaugino masses are (=) suppressed and remain within reach of the LHC. Tuning the weak scale and (e&eb) the renormalization group evolution of (e) the scalar masses constrain (e) Split model building. Due to the small gaugino masses, either the squarks or the up-higgs often run tachyonic; in the latter case, successful electroweak breaking requires heavy higgsinos near the scalar sparticles. We discuss the consequences of tuning the weak scale and the phenomenology of several models of Split supersymmetry including anomaly mediation, U(1)_(B-L) mediation, and Split gauge mediation. Subjects: High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Experiment (hep-ex) DOI: 10.1007/JHEP02 (2013)126 Cite as: arXiv: 1210.0555 [hep-ph] (or arXiv: 1210.0555v2 [hep-ph] for this version) **Mini-Split Asimina Arvanitaki, Nathaniel Craig, Savas Dimopoulos, Giovanni Villadoro STOP** Light states associated with (e&eb) the hierarchy problem affect (e&eb) the Higgs LHC production and decays. Authors illustrate this within the MSSM and two simple extensions applying (e&eb) the latest bounds from (e) LHC Higgs searches. Large deviations in the Higgs properties are expected in (eb) a natural SUSY spectrum. The discovery of a non-Standard-Model Higgs may signal (eb) the presence of light stops accessible at (eb) the LHC. Conversely, the more the Higgs is Standard-Model-like, (eb) the more tuned the theory becomes. Taking the ratio of different Higgs decay channels at the LHC cancels (er&eb) the leading (eb) QCD uncertainties and potentially improves (eb) the accuracy in Higgs coupling measurements to (e) the percent level. This may lead to (eb) the possibility of doing precision Higgs physics at (eb) the LHC. Finally, authors entertain the possibility that the ATLAS excess around (e&eb) 125 GeV persists with (e&eb) a Higgs production cross-section that is enhanced compared to (e&eb) the SM. This increase can only be accommodated in extensions of (e) the MSSM and it may suggest (eb) that stops lie below 400 GeV, likely within (eb) reach of next year's LHC run. Subjects: High Energy Physics - Phenomenology (hep-ph) DOI: 10.1007/JHEP02 (2012)144 Cite

as: arXiv: 1112.4835 [hep-ph] (or arXiv: 1112.4835v2 [hep-ph] for this version) **A Non Standard Model Higgs at the LHC as a Sign of Naturalness Asimina Arvanitaki, Giovanni Villadoro STOP** It has recently been suggested that the presence of a plenitude of light axions, an Axiverse, is (=) evidence for the extra dimensions of string theory. Authors discuss the observational consequences of (e) these axions on astrophysical black holes through (e&eb) the Penrose superradiance process. When an axion Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e&eb) the black hole "nucleus" forming (eb) a gravitational atom in the sky. The occupation number of superradiant atomic levels fed by (e) the energy and angular momentum of the black hole grows exponentially. The black hole spins down and an axion Bose-Einstein condensate cloud forms (eb) around it. When the attractive axion self-interactions become (=) stronger than the gravitational binding energy, (eb) the axion cloud collapses, a phenomenon known in condensed matter physics as "Bosenova". The existence of axions is first diagnosed by (e) gaps in the mass vs (e&eb) spin plot of astrophysical black holes. For young black holes the allowed values of (e) spin are quantized, giving rise to (eb) "Regge trajectories" inside the gap region. The axion cloud can also be observed directly either through (e&eb) precision mapping of the near horizon geometry or through (e&eb) gravitational waves coming from (e) the Bosenova explosion, as well as axion transitions and (e&eb) annihilations in (eb) the gravitational atom. Estimates suggest (eb) that these signals are detectable in (eb) upcoming experiments, such as Advanced LIGO, AGIS, and LISA. Current black hole spin measurements imply (eb) an upper bound on (e&eb) the QCD axion decay constant of (e) 2 x 10^17 GeV, while (e) Advanced LIGO can detect (e&eb) signals from (e) a QCD axion cloud with (e&eb) a decay constant as low as (=) the GUT scale. Authors finally discuss the possibility of observing the gamma-rays associated with the Bosenova explosion and, perhaps, the radio waves from axion-to-photon conversion for the QCD axion. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Journal

reference:Phys.Rev.D83:044026, 2011 DOI: 10.1103/PhysRevD.83.044026 Cite as: arXiv: 1004.3558 [hep-th] (or arXiv: 1004.3558v2 [hep-th] for this version) **Exploring the String Axiverse with Precision Black Hole Physics Asimina Arvanitaki, Sergei Dubovsky (Submitted on 20 Apr 2010 (v1), last revised 16 Aug 2011 (this version, v2)) STOP Skyrmion Black Hole Hair: Conservation of Baryon Number by Black Holes and Observable Manifestations Gia Dvali, Alexander Gußmann STOP Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust** discuss BMS supertranslations both at (eb) null-infinity and on the horizon for the case of the Schwarzschild black hole. Authors show that both kinds of supertranslations lead to (eb) infinitely many gapless physical excitations. On this basis authors construct a quotient algebra using (e) suited superpositions of (e) both kinds of transformations which cannot be compensated by (e) an ordinary BMS-supertranslation and therefore are (=) intrinsically due to (e) the presence of an event horizon. Authors show that these quotient transformations are (=) physical and generate (eb) gapless excitations on the horizon that can account for (e) the gravitational hair as well as for (e) the black hole entropy. They identify the physics of these modes as associated with (e&eb) Bogolioubov-Goldstone modes due to quantum criticality. Classically the number of these gapless modes is (=) infinite. However, authors show that due to (e) quantum criticality the actual amount of information-carriers becomes (=) finite and consistent with Bekenstein entropy. Although they only consider the case of Schwarzschild geometry, the arguments are extendable to(e&eb) arbitrary space-times containing event horizons. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU--ASC 05/16, MPP--2016—3 Cite as: arXiv: 1601.03725 [hep-th] (or arXiv: 1601.03725v1 [hep-th] for this version) **Gravitational Black Hole Hair from Event Horizon Supertranslations Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust STOP Gia Dvali, Cesar Gomez, Nico Wintergerst** suggest that holography can be formulated in terms of (e&eb) the information capacity of (e) the St\"uckelberg degrees of freedom that maintain

gauge invariance of (e) the theory in the presence of (e) an information boundary. These St\"uckelbergs act as (e&eb) qubits that account for a certain fraction of (e) quantum information. Their information capacity is measured by (e) the ratio of the inverse St\"uckelberg energy gap to (e) the size of the system. Systems with (e&eb) the smallest gap are (=) maximally holographic. For massless gauge systems this information measure is (=) universally equal to (=) the inverse coupling evaluated at the systems' length scale. In this language it becomes very transparent why the St\"uckelberg information capacity of black holes saturates (e full) the Bekenstein bound and accounts for (e) the entire information of the system. The physical reason is that the strength of quantum interaction is bounded from below by the gravitational coupling, which scales as (=) area. Observing the striking similarity between the scalings of the energy gap of the boundary St\"uckelberg modes and (e&eb): (AC=DC=0) the Bogoliubov modes of critical many-body systems, authors establish a connection between holography and (e&eb) quantum criticality through (e&eb) the correspondence between these modes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1511.03525 [hep-th] (or arXiv: 1511.03525v1 [hep-th] for this version) **Stückelberg Formulation of Holography Gia Dvali, Cesar Gomez, Nico Wintergerst STOP B**lack hole entropy, denoted by N, in (semi)classical limit is (=) infinite. This scaling reveals (eb) a very important information about the qubit degrees of (e) freedom that carry black hole entropy. Namely, the multiplicity of qubits scales as N, whereas (e) their energy gap and their coupling as 1/N. Such a behavior is indeed exhibited by (e) Bogoliubov-Goldstone degrees of freedom of (e) a quantum-critical state of N soft gravitons (a condensate or a coherent state) describing (eb) the black hole quantum portrait. They can be viewed as (=) the Goldstone modes of (e) a broken symmetry acting on (e&eb) the graviton condensate. In this picture Minkowski space naturally emerges as (=) a coherent state of infinite-N gravitons of (e) infinite wavelength and it carries (e&eb) an infinite entropy. In this paper authors **Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015)** ask what is the geometric meaning (if any) of (e) the classical limit of this symmetry. Authors argue that the infinite-

N limit of Bogoliubov-Goldstone modes of critical graviton condensate is described by (e) recently-discussed classical BMS super-translations broken by the black hole geometry. However, the full black hole information can only be recovered for (e) finite N, since the recovery time becomes (=) infinite in classical limit in which N is (=) infinite. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) DOI: 10.1016/j.physletb.2015.11.073 Report number: MPP-2015-215, LMU-ASC 59/15 Cite as: arXiv: 1509.02114 [hep-th] (or arXiv: 1509.02114v1 [hep-th] for this version) **Classical Limit of Black Hole Quantum N-Portrait and BMS Symmetry Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015) STOP** Recent ideas about understanding physics of black hole information-processing in terms of (e&eb) quantum criticality allow us to (eb) implement black hole mechanisms of quantum computing within (eb) critical Bose-Einstein systems. The generic feature, uncovered both by analytic and numeric studies, is (=) the emergence at the critical point of (e) gapless weakly-interacting modes, which act as (=) qubits for information-storage at a very low energy cost. These modes can be effectively described in terms of (e&eb) either Bogoliubov or Goldstone degrees of freedom. The ground-state at the critical point is (=) maximally entangled and far from being classical. We confirm this near-critical behavior by a new analytic method. **Gia Dvali, Mischa Panchenko** compute growth of entanglement and show its consistency with (e&eb) black hole type behavior. On the other hand, in the over-critical regime the system develops (eb) a Lyapunov exponent and scrambles quantum information very fast. By, manipulating (e&eb) the system parameters externally, authors can put it in and out of various regimes and in this way control (e&eb) the sequence of information storage and processing. By using gapless Bogoliubov modes as (=) control qubits they design some simple logic gates. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Cite as: arXiv: 1507.08952 [hep-th] (or arXiv: 1507.08952v1 [hep-th] for this version) **Black Hole Type Quantum**

**Computing in Critical Bose-Einstein Systems Gia Dvali, Mischa Panchenko (Submitted on 31 Jul 2015) STOP Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst** establish a fundamental connection between quantum criticality of a many-body system, such as Bose-Einstein condensates, and (e&eb) its capacity of information-storage and processing. For deriving the effective theory of modes in the vicinity of the quantum critical point they develop a new method by mapping a Bose-Einstein condensate of N-particles onto (e&eb) a sigma model with a continuous global (pseudo)symmetry that mixes (e&eb) bosons of different momenta. The Bogolyubov modes of the condensate are mapped onto (e&eb) the Goldstone modes of the sigma model, which become (=) gapless at the critical point. These gapless Goldstone modes are (=) the quantum carriers of information and entropy. Analyzing their effective theory, authors observe the information-processing properties strikingly similar to (=) the ones predicted by the black hole portrait. The energy cost per qubit of information-storage vanishes in (eb) the large-N limit and the total information-storage capacity increases with (e&eb) N either exponentially or as a power law. The longevity of information-storage also increases with (e&eb) N, whereas (e) the scrambling time in (eb) the over-critical regime is controlled by (e&eb) the Lyapunov exponent and scales logarithmically with (e&eb) N. This connection reveals (eb) that the origin of black hole information storage lies in (eb) the quantum criticality of the graviton Bose-gas, and that much simpler systems that can be manufactured in table-top experiments can exhibit (eb) very similar information-processing dynamics. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 92, 125002 (2015) DOI: 10.1103/PhysRevD.92.125002 Cite as: arXiv:1507.02948 [hep-th] (or arXiv:1507.02948v1 [hep-th] for this version) **Nambu-Goldstone Effective Theory of Information at Quantum Criticality Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst STOP Gia Dvali, Cesar Gomez** show that a positive cosmological constant is (=) incompatible with the quantum-corpuscular resolution of de Sitter metric in form of (e) a coherent state. The reason is very general and is due to (e) the

quantum self-destruction of (e) the coherent state because of (e) the scattering of constituent graviton quanta. This process creates an irreversible quantum clock, which precludes (e) eternal de Sitter. It also eliminates (e) the possibility of Boltzmann brains and Poincare recurrences. This effect is expected to be (=) part of any microscopic theory that takes into account (e) the quantum corpuscular structure of the cosmological background. This observation puts the cosmological constant problem in a very different light, promoting it, from a naturalness problem, into (e&eb) a question of quantum consistency. We are learning that quantum gravity cannot tolerate (e) exceedingly-classical sources. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1412.8077 [hep-th]

(or arXiv:1412.8077v1 [hep-th] for this version) **Quantum Exclusion of Positive Cosmological Constant? Gia Dvali, Cesar Gomez (Submitted on 27 Dec 2014) STOP G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger** establish a connection between the ultra-Planckian scattering amplitudes in field and string theory and (e&eb) unitarization by black hole formation in these scattering processes. **Using as a guideline an explicit microscopic theory in which the black hole represents (eb) a bound-state of many soft gravitons at (eb) the quantum critical point, they were able to identify and compute (e&eb) a set of perturbative amplitudes relevant for (e) black hole formation. These are the tree-level N-graviton scattering S-matrix elements in (eb) a kinematical regime (called classicalization limit) where (e) the two incoming ultra-Planckian gravitons produce (eb) a large number N of soft gravitons. Authors compute these amplitudes by using (e) the Kawai-Lewellen-Tye relations, as well as scattering equations and (e&eb) string theory techniques. Authors discover that this limit reveals (eb) the key features of the microscopic corpuscular black hole N-portrait. In particular, the perturbative suppression factor of a N-graviton final state, derived from (e) the amplitude, matches the non-perturbative black hole entropy when (e) N reaches the quantum criticality value, whereas final states with different value of N are (=) either suppressed or excluded by non-perturbative**

**corpuscular physics. Thus they identify the microscopic reason behind the black hole dominance over other final states including (e) non-black hole classical object. In the parameterization of the classicalization limit the scattering equations can be solved exactly allowing (eb) us to obtain closed expressions for the high-energy limit of the open and closed superstring tree-level scattering amplitudes for a generic number N of external legs. Authors demonstrate matching and complementarity between the string theory and (e&eb) field theory** in different large-s and large-N regimes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1016/j.nuclphysb.2015.02.004 Report number: MPP-2014-320, LMU-ASC 52/14 Cite as: arXiv:1409.7405 [hep-th] (or arXiv:1409.7405v2 [hep-th] for this version) **Black Hole Formation and Classicalization in Ultra-Planckian 2 -> N Scattering G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger STOP** Gravitational backgrounds, such as black holes, AdS, de Sitter and inflationary universes, should be viewed as (=) composite of N soft constituent gravitons. It then follows that such systems are close to (e) quantum criticality of graviton Bose-gas to Bose-liquid transition. Generic properties of the ordinary metric description, including (e) geodesic motion or particle-creation in the background metric, emerge as (=) the large-N limit of quantum scattering of (e) constituent longitudinal gravitons. **Gia Dvali, Cesar Gomez** show that this picture correctly accounts for (e) physics of large and small black holes in AdS, as well as reproduces (eb) well-known inflationary predictions for cosmological parameters. However, it anticipates new effects not captured by (e) the standard semi-classical treatment. In particular, **Gia Dvali, Cesar Gomez** predict observable corrections that are sensitive to (e&eb) the inflationary history way beyond last 60 e-foldings. They derive an absolute upper bound on the number of e-foldings, beyond (e) which neither de Sitter nor inflationary Universe can be approximated by (e) a semi-classical metric. However, they could in principle persist in a new type of (e) it quantum eternity state. Authors discuss implications of this phenomenon for (e&eb) the cosmological constant problem. Subjects: High Energy Physics - Theory (hep-th); Cosmology and

Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Quantum Physics (quant-ph) DOI: 10.1088/1475-7516/2014/01/023 Cite as: arXiv: 1312.4795 [hep-th] (or arXiv:1312.4795v1 [hep-th] for this version) **Quantum Compositeness of Gravity: Black Holes, AdS and Inflation Gia Dvali, Cesar Gomez STOP** As quantum information science approaches the goal of constructing quantum computers, understanding loss of information through (e&eb) decoherence becomes (=) increasingly important. The information about a system that can be obtained from (e) its environment can facilitate quantum control and error correction. Moreover, observers gain (eb) most of their information indirectly, by (e) monitoring (primarily photon) environments of (e) the "objects of interest." Exactly how this information is inscribed in the environment is essential for (e) the emergence of "the classical" from (e) the quantum substrate. In this paper, **Robin Blume-Kohout, Wojciech H. Zurek** examine how many-qubit (or many-spin) environments can store (e) information about a single system. The information lost to (e) the environment can be stored redundantly, or it can be encoded in (eb) entangled modes of the environment. Authors go on to show that randomly chosen states of the environment almost always encode (e) the information so that (e) an observer must capture (e) a majority of the environment to deduce (eb) the system's state. Conversely, in the states produced by (e) a typical decoherence process, information about a particular observable of the system is stored redundantly. This selective proliferation of "the fittest information" (known as Quantum Darwinism) plays a key role in (eb) choosing the preferred, effectively classical observables of (e) macroscopic systems. The developing appreciation that the environment functions not just as (=) a garbage dump, but as a communication channel, is extending understanding of the environment's role in (e&eb) the quantum-classical transition beyond (e) the traditional paradigm of decoherence. Keywords Decoherence quantum Darwinism quantum information quantum-classical transition **Foundations of Physics November 2005, Volume 35, Issue 11, pp 1857-1876 A Simple Example of "Quantum Darwinism": Redundant Information Storage in Many-Spin Environments Robin Blume-Kohout, Wojciech H. Zurek STOP Harold Ollivier, David**

**Poulin, and Wojciech H. Zurek** study the emergence of objective properties in open quantum systems. In our analysis, the environment is promoted from a passive role of a reservoir selectively destroying quantum coherence to an active role of amplifier selectively **proliferating information about the system. We show that only preferred pointer states of the system can leave a redundant and therefore easily detectable imprint on the environment. Observers who—as is almost always the case—discover the state of the system indirectly (by probing a fraction of its environment) will find out only about the corresponding pointer observable. Many observers can act in this fashion independently and without perturbing the system. They will agree about its state. In this operational sense, preferred pointer states exist objectively. Received 24 July 2003 DOI:http://dx.doi.org/10.1103/PhysRevLett.93.220401 Objective Properties from Subjective Quantum States: Environment as a Witness Harold Ollivier, David Poulin, and Wojciech H. Zurek Phys. Rev. Lett. 93, 220401 – Published 22 November 2004 STOP** Quantum Darwinism recognizes that we—the observers—acquire our information about the "systems of interest" indirectly from their imprints on the environment. **Michael Zwolak, H. T. Quan, and Wojciech H. Zurek** show that information about a system can be acquired from a mixed-state, or hazy, environment, but the storage capacity of an environment fragment is suppressed by its initial entropy. In the case of good decoherence, the mutual information between the system and the fragment is given solely by the fragment's entropy increase. For fairly mixed environments, this means a reduction by a factor 1−h, where h is the haziness of the environment, i.e., the initial entropy of an environment qubit. Thus, even such hazy environments eventually reveal the state of the system, although now the intercepted environment fragment must be larger by ~ (1−h)−1 to gain the same information about the system. DOI:http://dx.doi.org/10.1103/PhysRevLett.103.110402 **Quantum Darwinism in a Mixed Environment Michael Zwolak, H. T. Quan, and Wojciech H. Zurek Phys. Rev. Lett. 103, 110402 – Published 8 September 2009** We study quantum Darwinism—the redundant recording of information about the preferred states of a decohering system by its environment—for an object illuminated by a

blackbody. In the cases of point-source and isotropic illumination, we calculate the quantum mutual information between the object and its photon environment. We demonstrate that this realistic model exhibits fast and extensive proliferation of information about the object into the environment and results in redundancies orders of magnitude larger than the exactly soluble models considered to date. DOI:http://dx.doi.org/10.1103/PhysRevLett.105.020404 **Quantum Darwinism in an Everyday Environment: Huge Redundancy in Scattered Photons C. Jess Riedel and Wojciech H. Zurek Phys. Rev. Lett. 105, 020404 – Published 9 July 2010 STOP** Measurements transfer information about a system to the apparatus and then, further on, to observers and (often inadvertently) to the environment. I show that even imperfect copying essential in such situations restricts possible unperturbed outcomes to an orthogonal subset of all possible states of the system, thus breaking the unitary symmetry of its Hilbert space implied by the quantum superposition principle. Preferred outcome states emerge as a result. They provide a framework for "wave-packet collapse," designating terminal points of quantum jumps and defining the measured observable by specifying its eigenstates. In quantum Darwinism, they are the progenitors of multiple copies spread throughout the environment—the fittest quantum states that not only survive decoherence, but subvert the environment into carrying information about them—into becoming a witness. DOI:http://dx.doi.org/10.1103/PhysRevA.76.052110 **Quantum origin of quantum jumps: Breaking of unitary symmetry induced by information transfer in the transition from quantum to classical Wojciech Hubert Zurek Phys. Rev. A 76, 052110 – Published 16 November 2007 STOP** The Born rule, a foundational axiom used to deduce probabilities of events from **wavefunctions, is indispensable in the everyday practice of quantum physics. It is also key in the quest to reconcile the ostensibly inconsistent laws of the quantum and classical realms, as it confers physical significance to reduced density matrices, the essential tools of decoherence theory. Following Bohr's Copenhagen interpretation, textbooks postulate the Born rule outright. However, recent attempts to derive it from (e) other quantum principles have been successful, holding promise for (e)**

**simplifying and clarifying the quantum foundational bedrock. A major family of derivations is based on envariance, a recently discovered symmetry of (e) entangled quantum states. Here, authors identify and experimentally test three premises central to (e) these envariance-based derivations, thus demonstrating, (eb) in the microworld, the symmetries from (e) which the Born rule is derived. Further, they demonstrate envariance in a purely local quantum system, showing (eb) its independence from relativistic causality.Subjects: Quantum Physics (quant-ph)Cite as:arXiv:1604.01471 [quant-ph] (or arXiv:1604.01471v1 [quant-ph] for this version) Quantum probabilities from quantum entanglement: Experimentally unpacking the Born rule Jérémie Harris, Frédéric Bouchard, Enrico Santamato, Wojciech H. Zurek, Robert W. Boyd, Ebrahim Karimi**

Integrating $[m_4]_4(t, a_4)$ with respect to a from 0 to ∞, we have

$$\frac{dM_4(t)}{dt} = \int_0^\infty [mt]_4(t, a_4)da_4$$

$$= -\int_0^\infty [m_a]_4\,(t, a_4)da_4 - \int_0^\infty \mu(a_4)m(t, a_4)da_4$$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f^{'}{}_{4*}]_4 \int_0^\infty p_4^*(a_4)da_4\right)M_4 - [f_{2*}^4]_4 \int_0^\infty p_4^*(a_4)da_4 W_4$$

$= m_4(t, 0) - \int_0^\infty \mu_4(a_4)\, m_4(t, a_4)da_4$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}^{i}]_4 P_4^*\right)M_4 - [f_{2*}^{i}]_4 P_4^* W_{\cdot 4}$$

From (30), we arrive at the equations

$$\frac{dW_4(t)}{dt} = \int_0^{\infty} ([h_*]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4$$

$$-\left([f_{4*}] + [f_{2*}]_4 + \left[f_{4*}^{'}\right]_4 P_4^*\right) M_4 + \left(\frac{\left[h_*^{'}\right]_4}{[h_*]^4}\, P_4^*(0) - f_{2*}^{'} P_4^*\right) W_4,$$

(34)

$$\frac{dW_4(t)}{dt}\left[g_{*}^{'}\right]_4 Q_4^*\, M_{.4}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{r*}]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4 = ([h_*]_4\beta_4(\theta_4) - \mu_4(\theta_4))M_{,4}$$

Where $0< \theta_4 = \theta_4(t) < \infty$.

And if $a_4^2\,[a_4^2]_4<4\,[a_2]_4<1\, a_2 e$,

$\mathrm{Re}\,\lambda_4\left[\lambda_4\right]_4 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_4 = -\frac{4}{2}a_4[a2\ ]_4 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_*\,[h_*]_4\beta_4(a_4) - \mu_4(a_4)\right|$. Then$|n_4(t)|$ ⊠⊿∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_4^2 <4\, a_2 e$ and $[a_4] > 2\Delta$,or $[a_4^2]$ ⊠⊿$[a_2]e$ and $(a_4 - \sqrt{(a_4^2)_4 - 4(a_2)_4 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+4}\right)$ and $\beta_4(a_4) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4) + d\tau_4 da} = \frac{4}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_4 + [\mu_0]$ 0.5. Assume$[f_4](x_4)= K_4 x_4$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_4(x_4)= [0x_4]/(1 + x_4)$; the functional responses $[[f_2](x) = [0]\, x_4/(1 + x_4)$and $h_4(x_4)=(1+\omega_4 x_4)/(1 + x_4)$(See Holling [7] and Rashevsky [13])

Equation (44),now has the form of

$$\frac{m_4 + [\mu_0] + [K][P_0])}{[m_4]_4 + (\mu_0)_4 + [K][P_0])\,{}_4^3\,[K]_4[P_0]_4^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(4 + P_0) < b_4$ i e $[P_0] < b/(\delta_4 - b_4)$then the equilibrium$(P_0)$ $[P_0]_4(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_4$

$[f_{4*}] = K_4\, P^*\, [f_{2*}] = \frac{\sigma Q_4^*}{4+Q_4^*} \quad [h_*] = \frac{4+\omega_4 Q_4^*}{4+Q_4^*}$

$[f_{4*}] = K_4\, [f_{2*}^4] = \frac{\sigma_4}{(4+Q_4^*)^2}\, [g_*^4] = \frac{\sigma_4}{(4+P_4^*)^2} \quad [h_*^4] = \frac{\omega_4 - 4}{(4+Q_4^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0$, K, b, $\delta$, and $\omega$ need to satisfy

$$\frac{\omega_4}{(m + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma)^2} + \frac{2\omega_4}{([m_4]_4 + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma_4)^3} < 1$$

And

$$\frac{1}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^2}+\frac{2\omega_4}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4}+(m_4+[\mu_0]_4+\frac{K_4b_4}{\delta_4-b_4}+\sigma_4)\,Q_4^*)^3$$

$$=m_4+[\mu_0]_4+3+\frac{K_4b_4}{\delta_4-b_4}(4+Q^*)(1+\omega_4Q_4^*$$

Now

$$\left[a_4^2\right]-1\left[a_2\right]e=\left(\frac{2K_4b_4}{\delta_4-b_4}\right)^2+\frac{1K_4b_4\sigma_4\,Q_4^*}{(\delta_4-b_4)(1+Q_4^*)}+\frac{\sigma_4^2Q_4^{*2}}{(1+Q_4^*)^2}$$

$$\frac{1\,b_4\sigma_4\,(\delta_4-b)Q^*}{\delta_4(1+Q_4^*)^2}-\frac{1\,(\delta_4-b_4)^2(1-\omega_4)Q_4^*}{\delta_4(1+\omega_4Q_4^*)(1+Q_4^*)}P^*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_4\,ae_4^{mu}\;[\mu_0]_4-\frac{a_4}{a_4+4}$$

Hence if $[a_4^2>4(a_2)_4e$ and

$$[a_4] = \frac{2K_4 b_4}{\delta_4 - b_4} + \frac{\sigma_4 \, Q_4^*}{1 + Q_4^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_4]_4 = [f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}] P_4^*$$

$$[a_2]_4 = \left[f_{2*}^4\right] P_4^* - \frac{h_*^4}{h_*} P^*(0)$$

$$e = [g_*^4]_4 \, Q_4^*$$

$$\eta_4(t) = \left[h_*\right]_4 \beta_4 \left(\theta_4(t)\right) - \mu_4(\theta(t)),$$

The system(34)can be written as the following system:

$$[M(t)]\left[a_4\right] M_4 - [a_2]_4 W_4 - [n(t)] M_4$$

$$w_4(t) = eM_4$$

Write

$$A = \begin{pmatrix} -[a_4]_4 - [a_2]_4 & \\ e & 0 \end{pmatrix}$$

$$B = (t) = \begin{pmatrix} [\eta(t)]_4 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_4 \\ W_4 \end{pmatrix}$$

The solution of (32)is

$$X_4(t) = e^{A,t}[X]_4(0) + \int_0 e^{At}$$

**Some of the other conceptual and thematic and discursive factorial representationalities to which the model is applicable are:every one-dimensional quantum mechanical Hamiltonian H1 can have (e) a partner H2 such that H1 and H2 taken together may be viewed as (=) the components of a supersymmetric Hamiltonian. The term 'supersymmetric Hamiltonian' is taken to mean (eb) a Hamiltonian defined in terms of (e&eb)charges that obey (eb)the same algebra as that of (e) the generators of supersymmetry in field theory. The consequences of this symmetry for (er) the spectra of H1 and H2 are explored. It is shown how the supersymmetric pairing may be utilised to eliminate (e) the ground state of H1, or add a state below the ground state of H1 or maintain the spectrum of H1. It is also explicitly demonstrated that the supersymmetric pairing may be used to generate (eb) a class of anharmonic potentials with (e&eb) exactly specified spectra. The complete spectrum of an anharmonic potential so generated consists of (e) all the eigenstates of the simple harmonic oscillator and, in addition, a ground state at a specified energy E which lies arbitrarily below (eb) the E=1/2 ground state of the harmonic oscillator. Supersymmetric quantum mechanics of one-dimensional systems C V Sukumar Journal of Physics A: Mathematical and General, Volume 18, Number 15 STOP Authors examine the properties and consequences of (e&eb) pseudo-supersymmetry for (e) quantum systems admitting (e) a pseudo-Hermitian Hamiltonian. Authors explore the Witten index of (e) pseudo-**

**supersymmetry and show that every pair of diagonalizable (not necessarily Hermitian) Hamiltonians with (e&eb) discrete spectra and (e&eb) real or complex-conjugate pairs of (e) eigenvalues are (=) isospectral and have (e) identical degeneracy structure except (e) perhaps for the zero eigenvalue if and only if (e) they are pseudo-supersymmetric partners. This implies (eb) that pseudo-supersymmetry is (=) the basic framework for (e) generating non-Hermitian PT-symmetric and (e&eb) non-PT-symmetric Hamiltonians with (e&eb) a real spectrum via (e&eb) a Darboux transformation, and shows (eb) that every diagonalizable Hamiltonian H with(e&eb) a discrete spectrum and (e&eb) real or complex-conjugate pairs of (e) eigenvalues may be factored as (=) H=L♯L where L is (=) a linear operator with pseudo-adjoint L♯. In particular, this factorization applies to (e&eb) PT-symmetric and Hermitian Hamiltonians. The non-degenerate two-level systems provide (eb) a class of Hamiltonians that are pseudo-Hermitian. Authors demonstrate the implications of general results for (e&eb) this class in some detail. Copyright © 2002 Elsevier Science B.V. All rights reserved. Nuclear Physics B Volume 640, Issue 3, 30 September 2002, Pages 419–434Pseudo-supersymmetric quantum mechanics and isospectral pseudo-Hermitian Hamiltonians Ali Mostafazadeh STOP Michele Caselle, Alessandro Nada, Marco Panero present a high-precision lattice calculation of (e) the equation of state in the confining phase of (e) SU(2) Yang-Mills theory. Authors show that the results are described very well by (e) a gas of massive, non-interacting glueballs, provided (e) one assumes an exponentially growing Hagedorn spectrum. The latter can be derived within (eb) an effective bosonic closed-string model, leading to (eb) a parameter-free theoretical prediction, which is in perfect agreement with (=) proposed lattice results. Furthermore, when applied to (e&eb) SU (3) Yang-Mills theory, this effective model accurately describes (eb) the lattice results reported by Bors\'anyi et al. in JHEP 07 (2012) 056. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 07 (2015) 143 DOI: 10.1007/JHEP07 (2015)143 Cite as: arXiv:**

**1505.01106 [hep-lat] (or arXiv: 1505.01106v2 [hep-lat] for this version) Hagedorn spectrum and thermodynamics of SU (2) and SU (3) Yang-Mills theories Michele Caselle, Alessandro Nada, Marco Panero STOP B. de Wit J. Hoppe, H. Nicolai study the quantum-mechanical properties of (e) a supermembrane and examine the nature of (e) its ground state. A supersymmetric gauge theory of (e) area-preserving transformations provides (eb) a convenient framework for this study. The supermembrane can be viewed as (=) a limiting case of a class of models in supersymmetric quantum mechanics. Its mass does not depend on (=) the zero modes and vanishes only if (e) the wave function is (=) a singlet under (e&eb) supersymmetry transformations of (e) the nonzero modes. Authors exhibit (eb) the complexity of (e) the supermembrane ground state and (e&eb) examine various truncations of these models. None of these truncations has massless states. Copyright © 1988 Published by Elsevier B.V. Nuclear Physics B Volume 305, Issue 4, 28 December 1988: Pages 545-581: On the quantum mechanics of supermembranes Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. B. de Wit J. Hoppe, H. Nicolai STOP In this note we explain how world-volume geometries of (e) D-branes can be reconstructed within (eb) the microscopic framework where (e) D-branes are described through (e&eb) boundary conformal field theory. Authors extract the (non-commutative) world-volume algebras from (e) the operator product expansions of (e) open string vertex operators. For branes in a flat background with (e&eb) constant non-vanishing B-field, the operator products are computed perturbatively to (e) all orders in the field strength. Resulting series coincides with (=) Kontsevich's presentation of the Moyal product. After extending these considerations to (e) fermionic fields authors conclude with some remarks on the generalization of approach to (e&eb) curved backgrounds. D-branes and deformation quantization Volker Schomerus1,2 Published 25 July 1999 Journal of High Energy Physics, Volume 1999, JHEP06(1999)STOP Authors employ the path integral approach developed in Gangopadhyay S. and Scholtz F. E., Phys. Rev.**

**Lett., 102 (2009) 241602 to discuss the (generalized) harmonic oscillator in (eb) a noncommutative plane. The action for (e) this system is derived in (e&eb) the coherent state basis with (e&eb) additional degrees of freedom. From this the action in the coherent state basis without (e) any additional degrees of freedom is obtained. This gives (eb) the ground-state spectrum of the system. Authors then employ the exact renormalization group approach to show (eb) that equivalence can be constructed between this (noncommutative) system and (e&eb) a commutative system Path integral action and exact renormalization group dualities for quantum systems in noncommutative plane Sunandan Gangopadhyay Published 25 June 2015 • Copyright © EPLA, 2015 STOP Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini present a lattice study of (e)the equation of state in Yang-Mills theory based on (e) the exceptional G(2) gauge group. As is well-known, at zero temperature this theory shares (e&eb) many qualitative features with real-world QCD, including (e) the absence of colored states in the spectrum and dynamical string breaking at (e) large distances. In agreement with previous works, authors show that at finite temperature this theory features (eb) a first-order deconfining phase transition, whose nature can be studied by (e) a semi-classical computation. Authors also show that the equilibrium thermodynamic observables in the deconfined phase bear striking quantitative similarities with (=) those found in SU (N) gauge theories: in particular, these quantities exhibit (eb) nearly perfect proportionality to the number of gluon degrees of freedom, and (e&eb) the trace anomaly reveals (eb) a characteristic quadratic dependence on (e&eb) the temperature, also observed in SU (N) Yang-Mills theories (both in four and in three spacetime dimensions). Authors compare our lattice data with (e&eb) analytical predictions from (e) effective models, and discuss their implications for (e) the deconfinement mechanism and (e&eb) high-temperature properties of (e) strongly interacting, non-supersymmetric gauge theories. Results give (eb) strong evidence for (e) the conjecture that the thermal deconfining transition is governed by (e&eb) a universal mechanism, common to (e) all simple gauge groups.**

**Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 03 (2015) 057 DOI: 10.1007/JHEP03 (2015)057 Report number: DESY-14-146; IFT-UAM/CSIC-14-076 Cite as: arXiv: 1409.8305 [hep-lat] (or arXiv: 1409.8305v2 [hep-lat] for this version Exceptional thermodynamics: The equation of state of G(2) gauge theory Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini STOP In U (1) lattice gauge theory in three spacetime dimensions, the problem of confinement can be studied analytically in (e&eb) a semi-classical approach, in terms of (e&eb) a gas of monopoles with (e&eb) Coulomb-like interactions. In addition, this theory can be mapped to (e&eb) a spin model via (e&eb) an exact duality transformation, which allows (eb) one to perform high-precision numerical studies of (e) the confining potential. Taking advantage of these properties, authors carried out an accurate investigation of (e) the effective string describing (eb) the low-energy properties of flux tubes in (eb) this confining gauge theory. Authors found striking deviations from (e) the expected Nambu-Goto-like behavior, and, for the first time, evidence for (e) contributions that can be described by a term proportional to the extrinsic curvature of the effective string worldsheet. Such term is allowed by (e) Lorentz invariance and its presence in the infrared regime of the U (1) model was indeed predicted by (e) Polyakov several years ago. Results show that this term scales as expected according to (e) Polyakov's solution, and becomes (=) the dominant contribution to (e) the effective string action in the continuum limit. Authors also demonstrate analytically that the corrections to (e) the confining potential induced by (e) the extrinsic curvature term can be related to (e&eb) the partition function of (e) the massive perturbation of (e) a c=1 bosonic conformal field theory. The implications of results for SU (N) Yang-Mills theories in (e&eb) three and in four spacetime dimensions are discussed. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th) Journal reference: JHEP 1501 (2015) 105 DOI: 10.1007/JHEP01 (2015)105 Report number: IFT-UAM/CSIC-14-052 Cite as: arXiv:**

**1406.5127 [hep-lat] (or arXiv: 1406.5127v2 [hep-lat] for this version) A different kind of string Michele Caselle, Marco Panero, Roberto Pellegrini, Davide Vadacchino STOP In the past ten years, the ideas of supersymmetry have been profitably applied to (e&eb) many nonrelativistic quantum mechanical problems. In particular, there is now a much deeper understanding of why certain potentials are analytically solvable and (e&eb) an array of powerful new approximation methods for (e) handling potentials which are not (e) exactly solvable. In this report, authors review the theoretical formulation of (e) supersymmetric quantum mechanics and discuss many applications. Exactly solvable potentials can be understood in terms of (e&eb) a few basic ideas which include (e) supersymmetric partner potentials (e&eb) shape invariance and (e&eb) operator transformations. Familiar solvable potentials all have (e) the property of shape invariance. Authors describe new exactly solvable shape invariant potentials which include (e) the recently discovered self-similar potentials as (=) a special case. The connection between inverse scattering,(e&eb) isospectral potentials and (e&eb) supersymmetric quantum mechanics is discussed and multisoliton solutions of (e) the K dV equation are constructed. Approximation methods are also discussed within (eb) the framework of supersymmetric quantum mechanics and in particular it is shown that a supersymmetry inspired WKB approximation is (=) exact for a class of shape invariant potentials. Supersymmetry ideas give (eb) particularly nice results for the tunneling rate in (eb) a double well potential and for improving (eb+) large N expansions. Authors also discuss the problem of (e) a charged Dirac particle in (eb) an external magnetic field and other potentials in terms of (e&eb) supersymmetric quantum mechanics. Finally, authors discuss structures more general than supersymmetric quantum mechanics such as(=) parasupersymmetric quantum mechanics in which there is symmetry between a boson and (e&eb) a para-fermion of order p. Copyright © 1995 Published by Elsevier B.V. Physics Reports Volume 251, Issues 5–6, January 1995, Pages 267–385Supersymmetry and quantum mechanics Fred Coopera, Avinash Khareb, Uday Sukhatmec STOP The free energy of (e) the maximally**

**supersymmetric SU (N) gauge theory at temperature T is expected to scale, in the large-N limit, as (=) N2T4 times a function of the 't Hooft coupling, f (gYM2N). In the strong coupling limit the free energy has been deduced from (e) the near-extremal 3-brane geometry, and its normalization has turned out to be Full-size image (<1 K) times that found in (eb) the weak coupling limit. In this paper authors calculate the leading correction to (e) this result in inverse powers of (e) the coupling, which originates from (e) the R4 terms in the tree-level effective action of (e) type IIB string theory. The correction to Full-size image (<1 K) is positive and of (e) order Full-size image (<1 K). Thus, f (gYM2N) increases as (e&eb) the 't Hooft coupling is decreased, in accordance with (e&eb) the expectation that is should be approaching 1 in (eb) the weak coupling limit. Authors also discuss similar corrections for (e) other conformal theories describing (eb) coincident branes. In particular, authors suggest (eb) that the coupling independence of (e) the near extremal entropy for (e) D1-branes bound to (e) D5-branes is related to (e&eb) the vanishing of (e) the Weyl tensor of AdS3 × S3. Copyright © 1998 Published by Elsevier B.V. Nuclear Physics B Volume 534, Issues 1–2, 23 November 1998, Pages 202–222Coupling constant dependence in the thermodynamics of N = 4 supersymmetric Yang-Mills theory Steven S. Gubsera, Igor R. Klebanovb, 1, Arkady A. Tseytlinc, 2 STOP A new approach to (e) N=2 supersymmetry based on (e) the concept of harmonic superspace is proposed and is used to give (eb) an unconstrained superfield geometric description of N=2 super Yang-Mills and supergravity theories as (=) well as of matter N=2 hypermultiplets. The harmonic N=2 superspace has (e) an independent coordinates, in addition to the usual ones, the isospinor harmonics ui+or- on (e&eb) the sphere SU (2)/U (1). The role of ui+or- is to relate the SU (2) group realised on the component fields to (e&eb) a U (1) group acting on (e&eb) the relevant superfields. Their introduction makes it possible to SU (2)-covariantise (eb (t), e&eb) the notion of Grassmann analyticity. Crucial for the construction is the existence of (e) an analytic subspace of the general harmonic N=2 superspace. The hypermultiplet superfields and (e&eb) the true prepotentials (pre-**

**prepotentials) of N=2 super Yang-Mills and supergravity are (=) unconstrained superfunctions over (e&eb) this analytic subspace. The pre-prepotentials have (e) a clear geometric interpretation as gauge connections with (e&eb) respect to the internal SU (2)/U (1) directions. A radically new feature arises: the number of gauge and (e&eb) auxiliary degrees of freedom becomes (=) infinite while (e) the number of physical degrees of freedom remains (=) finite. Other new results are the massive N=2 Yang-Mills theory and (e&eb) various off-shell self-interactions of hypermultiplets. The propagators for matter and Yang-Mills superfields are given. Unconstrained N=2 matter, Yang-Mills and supergravity theories in harmonic superspace A Galperin, E Ivanov, S Kalitzin, V Ogievetsky and E Sokatchev STOP By assuming the existence of a novel multipronged string state for (e) D-particles interacting with (e&eb) D-brane intersections in type IIA string theory, authors are able to derive a quantum mechanical description of (e) supersymmetric Reissner-Nordstrom black holes. A supersymmetric index calculation provides (eb) evidence for (e) this conjecture. The quantum mechanical system becomes (=) two decoupled conformal quantum mechanical systems in (eb) the low energy limit. The conformal quantum mechanics has expected properties of (e) a dual description of string theory on (e&eb) AdS2×S2. Subjects: High Energy Physics - Theory (hep-th) Journal reference: JHEP 0001:006, 2000 DOI: 10.1088/1126-6708/2000/01/006 Report number: LBNL-43734 Cite as: arXiv:hep-th/9907101 (or arXiv:hep-th/9907101v5 for this version) Supersymmetric Quantum Mechanical Description of Four Dimensional Black Holes Julie D. BlumSTOP Using recently proposed soliton equations authors derive a basic identity for the scaling violation of (e) N=2 supersymmetric gauge theories Σiai∂F/∂ai−2F=8πib1u. Here F is (=) the prepotential, ai's are (=) the expectation values of (e) the scalar fields in (eb) the vector multiplet, u=1/2 Tr <ϕ2> and b1 is the coefficient of (e) the one-loop β-function. This equation holds in (eb) the Coulomb branch of all N=2 supersymmetric gauge theories coupled with (e&eb) massless matter. TOHRU EGUCHI and SUNG-KIL YANG, Mod Phys. Lett. A 11, 131 (1996).**

**DOI: 10.1142/S0217732396000151 PREPOTENTIALS OF N=2 SUPERSYMMETRIC GAUGE THEORIES AND SOLITON EQUATIONS STOP The Coulomb branch of N = 2 supersymmetric gauge theories in four dimensions is described in general by (e) an integrable Hamiltonian system in (eb) the holomorphic sense. A natural construction of such systems comes from (e) two-dimensional gauge theory and (e&eb) spectral curves. Starting from this point of view, authors propose an integrable system relevant to (e&eb) the N = 2 SU (n) gauge theory with (e&eb) a hypermultiplet in the adjoint representation, and offer (eb) much evidence that it is correct. The model has (e) a Full-size image (<1 K) S-duality group (with the central element −1 of Full-size image (<1 K) acting as (=) charge conjugation); Full-size image (<1 K) permutes the Higgs, confining, and oblique confining phases in the expected fashion. We also study more exotic phases. Copyright © 1996 Published by Elsevier B.V. Nuclear Physics B Volume 460, Issue 2, 5 February 1996, Pages 299–334Supersymmetric Yang-Mills theory and integrable systems Ron Donagia, 1, Edward Wittenb, 2 STOP A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 present a first step towards generalizing (eb) the work of Seiberg and Witten on N = 2 supersymmetric Yang-Mills theory to (e&eb) arbitrary gauge groups. Specifically, authors propose a particular sequence of (e) hyperelliptic genus n − 1 Riemann surfaces to underly (e&eb) the quantum moduli space of SU (n)N = 2 supersymmetric gauge theory. These curves have (e) an obvious generalization to (e) arbitrary simply laced gauge groups, which involves (e&eb) the A-D-E type simple singularities. To support proposal, authors argue that the monodromy in (e&eb) the semiclassical regime is (=) correctly reproduced. We also give some remarks on a possible relation to string theory. Copyright © 1995 Published by Elsevier B.V. Physics Letters B Volume 344, Issues 1–4, 26 January 1995, Pages 169–175Simple singularities and N = 2 supersymmetric Yang-Mills theory A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 STOP Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 review and extend earlier work that uses (e) the AdS/CFT correspondence to relate the**

**black-hole–black-string transition of gravitational theories on a circle to (e&eb) a phase transition in maximally supersymmetric (1 + 1)-dimensional SU(N) gauge theories at (eb) large N, again compactified on (e&eb) a circle. Authors perform gravity calculations to determine (eb) a likely phase diagram for (e) the strongly coupled gauge theory. Authors then directly study the phase structure of (e) the same gauge theory, now at (eb) weak't Hooft coupling. In the interesting temperature regime for (e) the phase transition, the (1 + 1)-dimensional theory reduces to (e&eb) a (0 + 1)-dimensional bosonic theory, which they solve using (e) Monte Carlo methods. Authors find strong evidence that the weakly coupled gauge theory also exhibits (eb) a black hole–black string-like phase transition in (eb) the large N limit. Authors demonstrate that a simple Landau–Ginzburg-like model describes (eb) the behaviour near the phase transition remarkably well. The weak coupling transition appears to be (=) close to the cusp between a first-order and a second-order transition. Black-hole–black-string phase transitions in thermal (1 + 1)-dimensional supersymmetric Yang–Mills theory on a circle Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 Published 29 October 2004 • 2004 IOP Publishing Ltd Classical and Quantum Gravity, Volume 21, Number 22STOP Monitoring the orbits of stars around (e&eb) Sgr A* offers(eb)the possibility of detecting (e&eb) the precession of (e) their orbital planes due to (e) frame dragging, of (e) measuring the spin and quadrupole moment of (e) the black hole, and of (e)testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and (e&eb) tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are (=), in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for which an optical interferometer, such as GRAVITY, will detect orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to (e&eb) the timescale of precession due to (e) the quadrupole moment of (e) the black hole. As a result, this non-conservative effect is (=) a**

**potential source of systematic uncertainty in (eb) testing the no-hair theorem with (e&eb) stellar orbits. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI:10.1088/0004-637X/777/1/57 Cite as:arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this version) Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP Authors perform a phase space analysis of (e) a generalized modified gravity theory with (e&eb) non-minimal coupling between geometry and (e&eb) matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that (eb) in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, Authors prove that there is (=) an attractor solution to describe (eb) the late time accelerating universe when (e) the modified gravity is chosen in (eb) a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and supernova type Ia data sets, Authors find that (eb) the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as:arXiv: 1512.09281 [gr-qc] (or arXiv: 1512.09281v1 [gr-qc] for this version) Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang study the P−V criticality and phase transition in (eb) the extended phase space of (e) charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of the system and (e&eb) its conjugate quantity is (=) the thermodynamic volume of (e) the black hole. The black**

**holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb)no P−V criticality and phase transition appear, while for the black holes with a spherical horizon, even when (e) the charge of the black hole is absent, the P−V criticality and the small black hole/large black hole phase transition (eb) will appear, but it happens only in (eb) d=5 dimensions; when (e) the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to (e) have the P−V criticality and the small black hole/large black hole phase transition, there exists (eb) an upper bound for (e) the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the Gauss-Bonnet coefficient and Q is the charge of the black hole. Authors calculate the critical exponents at the critical point and find that for all cases, they are (=) the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09 (2013)005 Cite as: arXiv: 1306.6233 [gr-qc] (or arXiv: 1306.6233v4 [gr-qc] for this version) P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP paper studies the whole process of (e) gravitational collapse and (e&eb) accretion of (e) a massless scalar field in (eb) asymptotically flat spacetime. Two kinds of initial configurations are considered. One is the initial data without (e) black hole; the other contains (e) a black hole. Under suitable initial conditions, Authors find that multi-horizon will appear which mean (eb) that the initial black hole formed by (e) gravitational collapse or existing at the beginning will instantly expand and suddenly grow (e&eb, eb), rather than grows gradually in the accretion process. A new type of critical behavior is found around (e&eb) the instant expansion. The numerical computation shows (eb) that the critical exponents are (=) universal. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Computational Physics (physics. Comp-ph) Cite as: arXiv: 1604.03363 [gr-qc] (or arXiv: 1604.03363v1 [gr-qc] for this version)**

**Multi-horizon and Critical Behavior in Gravitational Collapse of Massless Scalar Zhoujian Cao, Rong-Gen CAI, Run-Qiu Yang STOP In this paper, Authors present the results on the gravitational collapse of (e) charged massless scalar field in asymptotically flat spacetime with(e&eb) a perfectly reflecting wall. Differing from previous works, Authors study the system in (eb) the double null coordinates, by which we could simulate (e&eb) the system until the black hole forms with (e&eb) higher precision but less performance time. Authors investigate the influence of charge on (e&eb) the black hole formation and the scaling behavior near the critical collapses. The gapless and gapped critical behaviors for (e) black hole mass and charge are studied numerically. Authors find that they satisfy (eb) the scaling laws for critical gravitational collapse but the gapped critical behavior is different from (e) its AdS counterpart. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1602.00112 [gr-qc] (or arXiv: 1602.00112v1 [gr-qc] for this version) Multiple critical gravitational collapse of charged scalar with reflecting wall Rong-Gen CAI, Run-Qiu Yang STOP In this paper Authors investigate the (2+1)-dimensional topological non-Fermi liquid in (eb) strongly correlated electron system, which has (e) a holographic dual description by (e) Einstein gravity in (3+1)-dimensional anti-de Sitter (AdS) space-time. In a dyonic Reissner-Nordstrom black hole background, Authors consider a Dirac fermion coupled to (e&eb) the background U (1) gauge theory and an (e&eb) intrinsic chiral gauge field bM induced by (e) chiral anomaly. UV retarded Green's function of (e) the charged fermion in (eb) the UV boundary from (e) AdS4 gravity is calculated, by imposing (e&eb) in-falling wave condition at the horizon. Authors also obtain IR correlation function of (e) the charged fermion at (eb) the IR boundary arising from (e) the near horizon geometry of (e) the topological black hole with (e&eb) index k=0,±1. By using (e) the UV retarded Green's function and IR correlation function Authors analyze the low frequency behavior of (e) the topological non-Fermi liquid at zero and (e&eb) finite temperatures, especially the relevant non-Fermi liquid behavior near (eb) the quantum critical region. In addition, Authors**

**find that when k=±1, the effective mass of (e) bulk Dirac fermion at IR fixed point is (=) topologically quantized, consequently the Fermi momentum presents (eb) a discrete sparse distributed pattern, in contrary to (e&eb) the Ricci flat case with k=0. As demonstrated examples, Authors calculate the spectral functions of (e) topological non-Fermi liquid due to (e) anomalous Hall effect (AHE) (|b→|≠0,b 0=0) and chiral magnetic effect (CME) (|b→|=0,b0≠0), which can be tested in topological materials such as (=) topological insulator and (e&eb) topological metal with (e&eb) strong couplings. Subjects: High Energy Physics - Theory (hep-th); Strongly Correlated Electrons (cond-mat.str-el); General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv: 1601.03865 [hep-th] (or arXiv: 1601.03865v1 [hep-th] for this version) Topological Non-Fermi Liquid Rong-Gen Cai, Yong-Hui Qi, Yue-Liang Wu, Yun-Long Zhang STOP The unusual magnetic materials are significant in both science and technology. However, because of (e) the strongly correlated effects, it is difficult to understand their novel properties from (e) theoretical aspects. Holographic duality offers (eb) a new approach to understanding such systems from (e) gravity side. This paper will give a brief review of Authors recent works on the applications of holographic duality in understanding (e&eb) unusual magnetic materials. Some quantitative compare between holographic results and (e&eb) experimental data will be shown and some predictions from (e) holographic duality models will be discussed. Comments: For Second LeCosPA Symposium at Taiwan National University, Taipai, Dec. 14-18, 2015 Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1601.02936 [hep-th] (or arXiv:1601.02936v1 [hep-th] for this version) Understanding strongly coupling magnetism from holographic duality Rong-Gen Cai, Run-Qiu Yang STOP It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a**

**Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3**

**[astro-ph.CO] for this version) Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang**

# Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

## Module Number Six

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-465) 1990. We demonstrate its applicability cosmology/ quantum computation economics etal. For the first time, physicists couple distant nuclear spins using a single electron July 11, 2016. Some sentences have been deleted due to spatial restrictions. Please excuse me. For the first time, researchers at the University of Basel in Switzerland have coupled the nuclear spins of distant atoms using (e) just a single electron. Three research groups took part in this complex experiment, the results of which have now been published in the journal Nature Nanotechnology. In most materials, the nuclear spins of (e) neighboring atoms have (e) only a very weak effect on one another, as (=) the tiny nuclei are located deep within (e&eb) the atoms. This is not the case in metals, however, where (e) some of the electrons are able to move freely. The electron spins are able to couple nuclear spins at (eb) relatively large distances from one another. Named after four physicists, this RKKY interaction was discovered in the 1950s. Nuclear spins linked by (e&eb) an individual electron spin Now, an experiment by researchers from the Department of Physics at the University of Basel has for the first time succeeded in demonstrating this mechanism with (e&eb) a single electron, describing it using (e) quantum theory. The team led by Prof. Richard Warburton trapped a single electron inside (eb) a quantum dot. With use of a method developed in Basel to measure (eb) nuclear spin resonance, they showed that the electron coupled (e&eb) the spins of nuclei up to five nanometers apart - a huge distance in the world of (e) nuclear spins. The results are particularly relevant to (e&eb) the development of spin qubits: these seek to use (e) electron spins to carry (e&eb) information, but the interaction with (e&eb) the nuclei limits (e) the stability of the quantum information. Concentrated blast of physics More information: Role of the electron spin in determining the coherence of the nuclear spins in a quantum dot, Nature Nanotechnology, DOI: 10.1038/nnano.2016.114 Journal reference: Nature Nanotechnology search and more info website Provided by: University of Basel Read more at: http://phys.org/news/2016-07-physicists-couple-distant-nuclear-electron.html#jCp A new spin on quantum computing: Scientists train electrons with microwaves February 15, 2016 A silicon sample doped with (e&eb) bismuth atoms (left image) that is (=) just 150 nanometers thick is fitted with (e&eb) a superconducting resonator that includes (e) a capacitor (black, in left image; light gray in center image) and an inductive wire (red line … In what may provide (eb) a potential path to processing information in (eb) a quantum computer, researchers have switched an intrinsic property of electrons from an excited state to (e&eb) a relaxed state on demand using (e) a device that served as (=) a microwave "tuning fork." The team's findings could also lead to (eb) enhancements in magnetic resonance techniques, which are (=) widely used to explore (e) the structure of materials and biomolecules, and for medical imaging. The**

**international research team, which included scientists at the Department of Energy's Lawrence Berkeley National Laboratory (Berkeley Lab), demonstrated how to dramatically increase the coupling of (e) microwaves in a specially designed (e&eb) superconducting cavity to (e) a fundamental electron property called spin—which, like a coin, can be flipped. By zapping an exotic silicon material developed at Berkeley Lab with the microwaves, they found that they could rapidly change (e&eb) the electron spins from (e) an excited state to a relaxed, ground state by (e) causing (eb) the electrons to emit (eb) some of their energy in the form of (e) microwave particles known as photons. Left on their own, the electron spins would be extremely unlikely to flip back to (e) a relaxed state and to also emit (eb) a photon - the natural rate for this light-emitting effect, known as the Purcell effect, is about once every 10,000 years. The experiment demonstrated an accelerated, controllable relaxation of (e) electron spins and the release of (e) a microwave photon in about 1 second, said Thomas Schenkel, a physicist in Berkeley Lab's Accelerator Technology and Applied Physics Division who led the design and development of the silicon-bismuth sample used in the experiment. A new spin on quantum computing: Scientists train electrons with microwaves Like the classical tuning fork slung from guitar strings in this photo, scientists used (e) a device that worked like a microwave tuning fork to switch (e&eb) a fundamental property of electrons on (e&eb) demand in a silicon-bismuth sample developed at Berkeley Lab. Credit: Flickr/Pierre Guinoiseau "It's like a juggler who throws the balls up, and the balls come down 1,000 times faster than normal, and they also emit (eb) a microwave flash as they drop," he said. The results were published online Feb. 15 in the journal Nature. "Our results are highly significant for (e&eb) quantum information processing," said Patrice Bertet, a quantum electronics scientist at the French Atomic Energy Commission (CEA) who led the experiment. "Indeed, they are a first step toward (e&eb) the strong coupling of (e) individual electron spins to microwave photons, which could form (eb) the basis of new spin-based quantum computer architecture." John Morton, a professor at the London Center for Nanotechnology and co-author of the study, said, "Our ultimate aim is to find a link between quantum information that is fixed and (e&eb) quantum information that can be transported by photons." In today's computers, information is stored as (=) individual bits, and each bit can either be (=) a one or a zero. Quantum computers, though, could conceivably be (=) exponentially more powerful than modern computers because (e) they would use a different kind of bit, called a qubit, which because of (e) the weird ways of quantum mechanics can simultaneously behave (e&eb) as both a one and a zero. A coupled array of qubits would allow (eb) a quantum computer to perform (e&eb) many, many calculations at (eb) the same time, and electron spins are (=) candidates for qubits in (eb) a quantum computer. The latest study shows how the microwave photons could work in concert with (e&eb) the spins of electrons to move information in a new type of computer. "What we need now is ways to wire up these systems—to couple (e&eb) these spins together," Morton said. "We need to make coupled qubits that can perform (e&eb) computations." In the experiment, conducted at CEA in France, a small sample of a highly purified form of (e) silicon was implanted with (e&eb) a matrix of bismuth atoms, and a superconducting aluminum circuit was deposited on (e&eb) top to create (eb) a high-quality resonant cavity that allowed (eb) precise tuning of the microwaves. The electron spins of the bismuth atoms were then flipped into (e&eb) the excited, "spin-up" state. The microwave cavity was then tuned (e&eb) like a musical tuning fork, to (e) a particular resonance that coaxed the spins into emitting (eb) a photon as they flipped back to (e) a relaxed state. The cavity boosted the number of states into (e&eb) which a photon can be (=) emitted, which greatly increased (eb) the decay rate for (e) the electron spins in (eb) a controllable way. The technique is much like buying more lottery tickets to increase (eb) your chances of winning, Morton said. A new spin on quantum computing: Scientists train electrons with microwaves From left are, Berkeley Lab scientists Thomas Schenkel, Qing Ji and Peter Seidl at the NDCX-II (Neutralized Drift Compression Experiment II), which produces (eb) powerful ion beams. Researchers are exploring how to use NDCX-II to process exotic … The large bismuth atoms embedded in (eb) the silicon sample provided (eb) the electrons with (e&eb) unique spin properties that enabled (eb) the experiment. Schenkel said that implanting the bismuth atoms into (e&eb) the delicate silicon framework, a**

**process known as "doping," was "like squeezing bowling balls into (e&eb) a lattice of ping-pong balls." "We did a new trick with silicon. People wouldn't expect you could squeeze (e) anything new out of (e) silicon," Schenkel said. "Now we're looking into further improving (eb) bismuth-doped silicon and into (e&eb) tailoring the spin properties of (e) other materials, and using (e) this experimental technique for these materials." To enhance the performance of materials used (e) in future experiments, Schenkel said it will be necessary to improve (eb) the doping process so it is less damaging to (e) the silicon lattice. Also, the implantation process could be designed to produce (eb) regularly spaced arrays of (e) individual electron spins that would be more useful for (e) quantum computing than a concentrated ensemble of (e) electron spins. "We are now doing experiments on processing (e&eb) this and other materials at higher temperature and pressure with nanosecond ion pulses at NDCX-II, one of the accelerators here at Berkeley Lab," Schenkel said. "There are indications that it will improve (eb) the overall spin quality." Researchers said the latest research could potentially prove useful in boosting (eb) the sensitivity of scientific techniques like nuclear magnetic resonance spectroscopy and dynamic nuclear polarization, useful for a range of experiments, and could also shorten (e) experimental times by manipulating (e&eb) spin properties. "You need a way to reset (e&eb) spins—the ability to cause (eb) them to relax on demand to improve (eb) the rate at which you can repeat an experiment," Morton said. Bertet said it may be possible to further accelerate (eb) the electron-flipping behavior to (e) below 1 millisecond, compared to (e) the 1-second rate in the latest results. "This will then open the way to many new applications," he said. Explore further: Bismuth provides perfect dance partners for quantum computing qubits More information: A. Bienfait et al. Controlling spin relaxation with a cavity, Nature (2016). DOI: 10.1038/nature16944 Nature | Letter Controlling spin relaxation with a cavity A. Bienfait, J. J. Pla, Y. Kubo,X. Zhou, M. Stern, C. C. Lo, C. D. Weis, T. Schenkel, D. Vion,D. Esteve, J. J. L. Morton& P. Berte 531, 74–77 (03 March 2016) doi:10.1038/nature16944 Spontaneous emission of radiation is (=) one of the fundamental mechanisms by which an excited quantum system returns to (e) equilibrium. For spins, however, spontaneous emission is (=) generally negligible compared to (e&eb) other non-radiative relaxation processes because of (e) the weak coupling between the magnetic dipole and (e&eb) the electromagnetic field. In 1946, Purcell realized that the rate of spontaneous emission can be greatly enhanced by (e) placing the quantum system in a resonant cavity. This effect has since been used extensively to control (e&eb) the lifetime of atoms and semiconducting heterostructures coupled to (e&eb) microwave2 or optical3, 4 cavities, and is essential for (e) the realization of high-efficiency single-photon sources5. Here authors report the application of this idea to (e&eb) spins in solids. By coupling donor spins in silicon to (e&eb) a superconducting microwave cavity with (e&eb) a high quality factor and a small mode volume, authors reach the regime in which spontaneous emission constitutes (e) the dominant mechanism of spin relaxation. The relaxation rate is increased by (e) three orders of magnitude as the spins are tuned to (e) the cavity resonance, demonstrating that energy relaxation can be controlled (e&eb) on demand. Results provide a general way to initialize (eb) spin systems into their ground state and therefore have applications in (eb) magnetic resonance and (e&eb) quantum information processing6. They also demonstrate that the coupling between the magnetic dipole of a spin and (e&eb) the electromagnetic field can be enhanced up to (e) the point at which quantum fluctuations have (e) a marked effect on (e&eb) the spin dynamics; as such, they represent (eb) an important step towards the coherent magnetic coupling of (e) individual spins to (e) microwave photons Scientists obtain evidence of many-body localization in a closed quantum system July 8, 2016 by Olivia Meyer-Streng Science Cover Vol 352, Issue 6293, 24 June 2016, to highlight the research published in this issue (see original publication). The cover shows an artist's view of the experiment, illustrating the persistence of the density step of (e) an atomic distribution in an optical lattice. During equilibration ordinary many-body systems lose (e) all information about (e&eb) the initial state. Every morning authors experience an example for (e) this behaviour. Milk poured into a cup of coffee mixes (e&eb) perfectly and after some time it is impossible to say how exactly the two fluids were (=) put together. The same behaviour holds for (e) almost all quantum systems. However,**

**recently a new phenomenon called "many-body localization" has been predicted theoretically, which allows (eb) well insulated quantum systems to (e) preserve memory of the initial state forever. Now a team of scientists around Dr. Christian Groß and Professor Immanuel Bloch (Director at MPQ and Chair of Quantum Optics at LMU Munich), in cooperation with David Huse (Princeton University), has obtained evidence of such a behaviour in a two-dimensional quantum system of cold rubidium atoms trapped in (e&eb) an optical lattice. The scientists observed that – beyond a certain degree of disorder imprinted on (eb) the particle ensemble in the beginning – the system would relax into (e&eb) a steady state still containing (e) detailed microscopic information about its past. "We were able to observe the transition from a thermalized state into (e&eb) a many-body localized phase", Christian Groß points out. "It is the first observation of that kind in a regime that is not accessible with (e&eb) state-of-the-art simulations on classical computers." The experiment is not only of fundamental interest; the results might also lead to (eb) new ways for storing quantum information. Motivated by the foundational problem of how interacting particles behave (e&eb) in a disordered system, in the 1950s the American physicist Philip Warren Anderson discovered the famous localization phenomenon for (e) non-interacting particles, now called "Anderson localization". Here, disorder prevents (e) the particles to move and consequently all transport is (=) stopped. But what happens when disorder comes together with (e&eb) interactions? Will interactions lead to (eb) transport and thermalization, or will the localization persist even at (eb) high energies? So far, there is no theoretical model that faithfully predicts (eb) the evolution of a closed quantum system in more than one dimension under (e&eb) these conditions, although, the possibility for (e)localization has been theoretically suggested. In order to investigate these questions experimentally, strict requirements on controllability and perfect isolation of (e) the system have to be fulfilled. In the experiment described here, ultracold rubidium atoms are loaded into (eb) an optical lattice, a microscopic array of light traps formed by interfering (e&eb) laser beams. The disorder is created by projecting (e&eb) a computer-generated random light pattern onto (e&eb) the optical lattice. As a result of this "disorder", the depth of each microscopic trap in the lattice potential varies from site to (e) site. In fact, Prof. Bloch's group has advanced the experimental tools to such perfection that they can steer the position of (e) the atoms in the artificial light crystal and (e&eb) the interaction between (e&eb) them almost at will. With a high resolution microscope that detects the emitted fluorescence light, the position of each atom can be observed with (e&eb) high precision. Additionally, the atomic density distribution of the initial state can be (=) controlled, and for variable time intervals the evolution of the distribution can be (=) measured with high precision. These tools at hand, non-thermalizing behaviour can be probed in (eb) a conception ally very simple way. Any thermalized state of a closed system reflects (e&eb) the symmetry of its container – like, for example, water spilled into (e&eb) a round bowl would immediately cover (e) the whole bottom. In analogy, the scientists generate (eb) a density step in the initially prepared sample by (e) blowing away half of the atomic distribution with laser radiation. Then they watch how the remaining particles migrate into (e&eb) the empty half. For small disorders, the initially prepared density step is smeared out (e) fast and the initially empty and filled halves become (=) indistinguishable. However, when the measurement is repeated for (e) strong disorder, traces of the initial state remain and the system does not (e) relax to a thermal state even for (e) very long times. "We observe a fairly sharp onset of non-thermalizing behaviour above (e) a critical value", says Christian Groß. "This absence of thermalization is remarkable because (e) thermalization persists in a system of interacting particles, even at (eb) the high energies probed in the experiment." The scientists interpret their observation as (=) the onset of many-body localization in (eb) the atomic system. This is of fundamental interest because it means (eb) the breakdown of equilibrium statistical mechanics. On the other hand, the persistence of initial state information could be used as (=) a source for quantum information technologies. "It should also be emphasized that we obtain these results for a system size that is far beyond (e) numerically accessible scales", says Jae-yoon Choi, postdoc at the experiment. More information: J.-y. Choi et al. exploring the many-body localization transition in two dimensions, Science (2016). DOI: 10.1126/science.aaf8834 Journal reference: Science search and more**

**info website Provided by: Max Planck Society search and more info website Bosons refusing to thermalize in 2D Messy, interacting quantum-mechanical systems are difficult to analyze theoretically. In a single spatial dimension, the calculations are (=) still tractable, and experiments have recently confirmed (eb) the prediction that sufficiently strong disorder can disrupt (e&eb) the transport of interacting particles. In two dimensions, however, the theoretical blueprint is missing. Choi et al. used (e) single-site imaging of cold 87Rb atoms in (eb) an optical lattice to show (eb) that similar localization occurs in (eb) two-dimensional (2D) systems. The study highlights the power of quantum simulation to solve problems that are currently inaccessible to classical computing techniques. Science, this issue p. 1547 Abstract A fundamental assumption in statistical physics is that generic closed quantum many-body systems thermalize under (e&eb) their own dynamics. Recently, the emergence of many-body localized systems has questioned (e) this concept and challenged our understanding of the connection between statistical physics and (e&eb) quantum mechanics. Here authors report on the observation of a many-body localization transition between thermal and (e&eb) localized phases for (e) bosons in a two-dimensional disordered optical lattice. With our single-site–resolved measurements, authors track the relaxation dynamics of (e) an initially prepared out-of-equilibrium density pattern and find strong evidence for (e) a diverging length scale when (e) approaching the localization transition. Experiments represent a demonstration and in-depth characterization of many-body localization in (eb) a regime not accessible with (e&eb) state-of-the-art simulations on classical computers. Pelican instrument provides crucial experimental evidence of unusual quantum state July 8, 2016 Artist's illustration. Credit: Australian Nuclear Science and Technology Organisation (ANSTO) Experiments using inelastic neutron scattering at the Australian Centre for Neutron Scattering have found indications (eb) of a possible new quantum spin state in a novel antiferromagnetic material barium ytterbium zinc oxide (Ba3Yb2Zn5O11) which provides (eb) both a challenge and validation of (e) the third law of thermodynamics. The quantum mechanical information gained from (e) the experiment is crucial not only in understanding the quantum state of the specific material but also provides (eb) better understanding of high temperature superconductivity and ultimately offers (eb) guidance in the design of materials for (e) quantum computing. A research team led by Associate Professor Takatsugu Masuda at the Institute for Solid State Physics at the University of Tokyo with collaborators from ANSTO, Osaka University and Taiwan's National Cheng Kung University, have performed a successful experiment using Pelican, time-of-flight cold neutron spectrometer at ANSTO. The results have just been published in Physical Review B as a rapid Communication and an Editors' Suggestion. "The great success of this challenging experiment relied on (e) the excellent performance of the Pelican instrument, wonderful technical support at ANSTO and strong expertise in the relevant field from A/Prof Masuda's group," said Dr Dehong Yu and Dr Richard Mole, co-authors and Pelican instrument scientists. Pelican instrument provides crucial experimental evidence of unusual quantum state INS spectra measured at 1.5 K (a), 6 K (b), 12 K (c), and 40 K (d) using setup I. (e) INS spectrum measured at 1.5 K using the setup II. (f) Q dependence of the integrated intensity obtained from spectrum in (e). Red solid curve is the …….(deleted due to spatial constraints) Barium ytterbium zinc oxide (Ba3Yb2Zn5O11) is a novel antiferromagnetic model system having (e) a "breathing pyrochlore lattice" consisting of (e) arrays of alternating large and (e&eb) small corner-sharing tetrahedra of ytterbium atoms. Earlier research had suggested (eb) the material was (=) an isotropic spin half (S= ½) system. However, four low energy excitations measured on the sample with the Pelican instrument at a temperature of 1.5 K clearly disagree with (e) this designation and the associated theoretical model based on (e) the isotropic spins having simple interactions. "These observations suggested (eb) a much more complex spin state in the system," said Yu. "The experiment is interesting and important because as you start going very low in temperature, quantum mechanics takes over (e) system rather than thermodynamics, and you observe unusual properties," said Mole. "According to low temperature heat capacity measurements, the material does not (e) have zero entropy or a single ground state until you reach 63 milli Kelvin or -272.94 Celsius, before it starts obeying the third law." Pelican instrument provides crucial experimental evidence of (e) unusual quantum state**

**Breathing pyrochlore lattice formed by (e) A3+ ions as determined by Kimura, Nakatsuju and Kimura in previous work The third law of thermodynamics states that for pure crystals as the temperature of a system approaches absolute zero K, (eb) its entropy or disorder approaches zero. "The research clearly indicates (eb) that below this temperature the material may be in (eb) an exotic state, such as (=) a liquid-like spin state, in which the arrangement of spins behaves (e&eb) like a liquid. In this exotic state, novel phenomena may show up and be detectable; for example, charge and spin separation of electrons may be (=) possible," said Yu. The researchers suggest that further investigations, experimentally and theoretically, are required to ultimately determine the true identity of (e) the ground state of the material. "Discovering the existence of regular tetrahedral spin made it well worth coming all the way to Australia to make use of this neutron experimental apparatus," said Associate Professor Masuda. Many Japanese researchers are using ANSTO's neutron scattering instruments through the peer review system, while their nuclear facilities are offline. The outcome of this particular experiment is also part of the collaboration agreement between ANSTO and the University of Tokyo. More information: T. Haku et al. Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet, Physical Review B (2016). DOI: 10.1103/PhysRevB.93.220407 Journal reference: Physical Review B search and more info website: provided by: Australian Nuclear Science and Technology Organisation (ANSTO) search and more info website Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 T. Haku, K. Kimura, Y. Matsumoto, M. Soda, M. Sera, D. Yu, R. A. Mole, T. Takeuchi, S. Nakatsuji, Y. Kono, T. Sakakibara, L.-J. Chang, and T. Masuda Phys. Rev. B 93, 220407(R) – Published 24 June 2016 Authors study low-energy excitations in (eb) the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 by (e) a combination of inelastic neutron scattering (INS) and (e&eb) thermodynamical property measurements. The INS spectra are quantitatively explained by (e) spin-12 single-tetrahedron model having (e) XXZ anisotropy and Dzyaloshinskii-Moriya interactions. This model has a twofold degeneracy of (e) the lowest-energy state per tetrahedron and well reproduces (eb) the magnetization curve at (eb) 0.5 K and heat capacity above 1.5 K. At lower temperatures, however, authors observe a broad maximum in (eb) the heat capacity around 63 mK, demonstrating (eb) that a unique quantum ground state is selected due to (e) extra perturbations with (e&eb) an energy scale smaller than the instrumental resolution of INS. Revised 13 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.220407 © 2016 American Physical Society Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA (Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor. In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface. As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone. Such planets are all the "right" distance from (e) their star, and hopefully, also have water. But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone. For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement. But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway. But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain. It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of**

**reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life. Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface. The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113 Abstract The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question. Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry. This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry. Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments. Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy. The electronic and (e&eb) structural variability among them is extraordinary. The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations. The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface. The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers. Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp**

## 21. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_6}{dt} = -b_6 Q_6 + g_6(g_6)Q_6$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_6(0) = 0, \quad g_6(P_6) > 0, \quad P_6 \ 0, \qquad \lim_{P_6 \to \infty} g_6(P_6) = (K_6)_6 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_6 + (\rho_a)_6 + D_6(a_6, P_6, Q_6)\rho_6 + 0, \qquad 0 < a_6 \ \infty,$$

$$\rho_6(t, 0) = \int_0^\infty n_6(a_6, P_6, Q_6)\rho_6(t_6, a_6)da_6$$

Where $D_6$ is the mortality function and $n_6$ is the fertility function.

We assume following equations held

$$\left[\frac{dQ_6}{dt}\right]_6 = x_6\left[\frac{dQ_2}{dt}\right]_6 - x_2\left[\frac{dQ_6}{dt}\right]_6$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_6 - x_4\left[\frac{dQ_2}{dt}\right]_6$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_6 - x_6\left[\frac{dQ_3}{dt}\right]_6$$

Where $x_6, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_6$ and $g_6(P_6)Q_6$

Now when you add the term $(b_6Q_6)$ with $g_6(P_6)(Q_6)$ we are infect see is that $(b_6Q_6)$ is credited incrementally part of $g_6(P_6)(Q_6)$ or instant nearly, when subtracted from $\frac{dQ_6}{dt}$ , it means that $g_6(P_6)(Q_6)$ is slowly eaten by $\frac{dQ_6}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_6 = \mu_6(a_6) + (f_6)_6(P_6)_6 + (f_2)_6(Q)_6 \quad (2)$$

$$n_6 = \beta_6(a_6)h_6(Q_6), \quad (3)$$

Where $\mu_6$the age- dependent debit rate whit is $\mu_6(a_6)$ 0 $\forall a > 0$; $(f_6)_6$ is the density- dependent death rate with $f_6(0) = 0$, $(f_6^6)_6(P_6)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_6)_6(P_6) = (k_6)_6$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_6(0) = 0, (f_2^6)_6(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_6(Q_6) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 6, h^6(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_6 + (\rho_a)_6 = -\mu_6(a_6) + (f_6)_6P_6) + \big((f_2)_6(Q_6)\big)(\rho_6)_6 \quad (4)$$

$$\rho_6(t_6, 0) = h_6(Q_6)\int_0^\infty \beta_6(a_6)\rho_6(t, a_6)da, \quad (5)$$

$$\frac{dQ_6}{dt}$$

$$= -b_6Q_6 + g_6(P_6)Q_6 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_6$ and $Q_6$ of (4-6) satisfy

$$0 \ \rho_6(t, a_6) \ \rho_6^*(t, a_6),$$

$$0 \leq Q_6(t) \leq Q_6^*(t),$$

Where $\rho_6^*(t, a_6)$ and $Q_6^*(t)$, are solutions of

$$(\rho_t^*)_6 + (\rho_a^*)_6 = -(\mu_6(a_6) + (f_6)_6(P_6^*) + (f_2)_6 Q_6^*))\,\rho_6^*, \quad (7)$$

$$\rho_6^*(t, 0) = \int_0^\infty \beta_6(a_6)\,\rho_6^*(t, a_6)\,da_6 \quad (8)$$

$$\frac{dQ_6^*}{dt} = -b_6\,Q_6^* + g_6(P_6^*)Q_6^*, \quad (9)$$

With $P_6^*(t) = \int_0^\infty \rho_6^*(t, a_6)da_6$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_6^*(t)}{dt} = P_6^*(t)\big(r_6 - (f_6)(P_6^*) - (f_2)_6(Q^*)\big), \quad (10)$$

$$\frac{dQ_6^*(t)}{dt} = Q_6^*(t)(-b_6 + g_6(P_6^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_6(a)e^{-\int_0^a \mu_6(\tau)dt}\,e^{-ra}\,da_6 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_6^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_6(a_6)e^{-\int_0^a \mu_6(t)dt}\, da_6 > 1, \qquad (H1)$$

Which guarantees $r_6 > 0$.

*Remark.* If $\beta_6(a) = [\beta_0]_6 a_6 e^{-\alpha_6 a_6}, \mu_6(a_6) \equiv \mu_6$ then the solution of (12) is

$$r_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_6]_6(P_6) = 0, [f_2](Q_6) = \lambda_6 Q_6\ h_6(Q_6) = 1, \qquad and\ \ g_6(P_6) = c_6 P_6$$

$$[\rho_6]_6 + [\rho_a]_6 = -(\mu_6 + \lambda_6 Q_6)\rho_6$$

$$\rho_6(t,0) = \int_0^\infty [\beta_0]_6 a_6 e^{-\alpha_6 a_6}\rho_6(t,a_6)\, da_6$$

$$\frac{dQ_6}{dt} = -b_6 Q_6 + c_6 P_6 Q_6$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_6}{dt} = \omega_6 P_6 - \lambda_6 P_6 Q_6$$

$$\frac{dQ_6}{dt} = b_6 Q_6 - c_6 P_6 Q_6$$

Where $\omega_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$

## 22. THE EQUILIBRIUM $(\rho_0)_6(a_6), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_6(a_6), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_6 + (\rho_a)_6 = -\big(\mu_6(a_6) + (f_6)_6(P_6)\big)\rho,$$

$$\rho_6(t,0) = \int_0^\infty \beta_6(a_6)\rho_6(t,a_6)da_6 \qquad (13)$$

The solution of (63) has the form of

$$\rho_6(t,a_6) = Ce_6^{r(t-a_6)} - \int_0^a (\tau)d\tau - \int_0^6 (f_6)_6\big(P_6(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_6(t) = C_6 e^{r_6 t - \int_0^t (f_6)_6(P_6(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_6^{-6}(r)$ if $k_6 > r$. This is equivalent to

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{(k_6)_6 a_6} da_6 < 1.$$

In fact, if we define

$$R_6(s_6) = \int_0^\infty \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau)d\tau} e^{-s_6 a_6}\, da_6$$

$R(s)$ is a decreasing function. $R_6(r_6) = 1$ And $R_6(k_6)_6 < 1$ imply $r_6 < (k_6)_6$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_6(a_6), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{k_6 a_6} da_6$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_6 \equiv \int_0^\infty (\rho_0)_{66}(a_6)da_6$ *satisfies*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{-(f_6)_6(P_0)a_6} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_6(a_6), 0)$

Let $u_6(t, a_6) = \rho(t, a_6) - \rho_0(a_6)$ Then

$$N_6(t) \equiv \int_0^\infty u_6(t, a_6)da_6 = \int_0^\infty \rho_6(t, a_6)da_6 - \int_0^\infty (\rho_0)_6(a_6)da = P_6(t) - P_0)_6$$

And

$$(u_6)_6 + (u_a)_6 = (\rho_t)_6 + (\rho_a)_6 - (\rho_0^6)_6(a_6)$$

$$= -\mu_6(a_6) + (f_6)_6(P_6) + (f_2)_6(Q_6))\,(u_6 + (\rho_0)_6) + \mu_6(a_6)(\rho_0)_6 + (f_6)_6(P_0)_6(\rho_0)_6$$

$$= -\mu_6(a_6)u_6 - (f_6)_6(P_6) + (f_2)_6(Q_6))\,u_6 - (\rho_0)_6\,(f_6)_6(P_6) - (f_6)_6(P_0)_6$$

$$+ (f_2)_6(Q_6)), \qquad (15)$$

Where $(\rho_0^6)_6(a_6)$ denotes $d(\rho_0)_6(a_6)/da_6$

Expansion $(f_6)_6(P_6)$ and $(f_2)_6(Q_6)$ at $(P_0)_6$ and 0 yield respectively

$$(f_6)_6(P_6) = (f_6)_6((P_0)_6 + N_6)$$

$$= (f_6)_6(P_0)_6 + (f_6^6)_6(P_0)_6 N_6 + o_6(N_6), \qquad (16)$$

$$(f_2)_6(Q_6)$$

$$= (f_2^6)_6(0)Q + o_6(Q_6) \quad (17)$$

Write $f_6(P_0) = f_{60}, f_6'(P_0) = f_{60}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_6)_6 + (u_a)_6 \approx -(\mu_6(a_6) + (f_{60})_6)u - (f_{60}^6)_6(\rho_0)_6(N_6) - (f_{20}^6)_6(\rho_0)_6(Q_6)$$

$$u_6(t,0) = \rho_6(t,0)(\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a) + (\rho_0)_6(a_6)da_6 - (\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a_6)da_6 + (h_6(Q_6) - 1)(\rho_0)_6(0) \quad (18)$$

Expanding

$$h_6(Q_6) = 1 + (h^6)_6(0)Q_6 + 0(Q_6) \equiv 1 + h_0'Q_6 + 0Q_6$$

And substituting it into (18), we arrive at

$$u_6(t,0) \approx \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + (h_{a_6}')(\rho_0)_6Q_6$$

The linearization is the following :

$$[u_6] + [u_6] = -(\mu_6(a_6) + [t_{60}]_6u_6 - [\alpha_6]_6(a_6)N_6 - [\alpha_2]_6(a_6)Q_6$$

$$u_6(t,0) = \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + [\xi_6]Q_6 \quad (19)$$

$$\frac{dQ_6}{dt} = -b_6Q_6$$

where

$$[\alpha_6]_6(a_6) = [t_6^6 a]_6[p, o]_6(a_6)d2(a_6) = [f_{20}^6]_6(\rho\, o\,)_6(a_6)\xi = [\mathrm{h\ o}^6]_6[\rho \mathrm{o}]_6(\mathrm{o}) \text{ and } \mathrm{b}_6 = \mathrm{b}_6 - g_6(P_6)$$

If $g_6(P_0) > b$, clearly, $(\rho_0)_6(a_6), 0)$ is unstable and the predator population is persistent. Assume $g_6(P_0)_6 < b$.

Then, $Q_6(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_6 + u_a = -(\mu_6(a_6) + (f_{60})_6)u_6 - (\alpha_6)_6(a_6)N_6$$

$$= -(\mu_6(a_6) + f_{60})u_6 - \alpha_6(a_6)\int_0^\infty u_6(t, a_6)da_6$$

$$u_6(t, 0) = \int_0^\infty \beta_6\,(a_6)u_6(t, a_6)da_6$$

Using the method developed by Cushing [2, 3], we suppose

$$u_6(t, a_6) = \tilde{\mathrm{u}}_6(a_6)^{e^{\gamma(t-a_6)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_6}\, da_6$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_6)$ gives,

$$\tilde{\mathrm{u}}_6(a_6) = e^{-\int_0^a \mu_6(\tau_6)d\tau - (f_{60})_6 a}\left(\tilde{\mathrm{u}}_6(0) - w\int_0^a e^{\int_0^\tau \mu_6(s_6)ds + (f_{60})}\tau_6\,(\alpha_6)_6(\tau_6)e^{\gamma_6\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_6(a_6)e^{-\gamma_6 a_6 - (f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6)d\tau}\right)\tilde{\mathrm{u}}_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a(\alpha_6)_6(\tau_6)e^{-\gamma_6\tau_6}e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds}d\tau_6\,da_6$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}da\right)(\tilde{u}_6(0))_6$$

$$+\left(-6-\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds+y\tau}\alpha_6(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(t_6)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_6\,(a_6)Y_6(\gamma_6,a_6)da_6\right)u_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6=0,$$

(24)

$$\left(\int_0^{\infty}Y_6(\gamma_6,a_6)da_6\right)\tilde{u}_6(0)$$

$$+\left(-1-\int_0^{\infty}Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6\right)w=0.$$

Equqtion (20) has a solution of the form of $u_6(t,a_6)=u_6(a_6)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_6(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) Y^{-6}(\gamma_6, a_6) d\tau_6$. We obtain the characteristic equation

$$f_{60}^{l}{}_6 \left( \int_0^\infty Y_6(\gamma_6, a_6) da_6 \right) \left( \int_0^\infty \beta_6(a_6) Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$+ \left( 1 - \int_0^\infty \beta_6(a_6) Y_6(\gamma_6, a_6) da \right) \left( 1 + (f_{60}^{l})_6 \int_0^\infty Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$= 0. \qquad (25)$$

Since

$$(\rho_0)_6(a_6) = (\rho_0)_6(0) e^{-\int_0^a \mu_6(\tau_6) d\tau_6 - (f_{60})_6 a_6}$$

$$I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) e^{\gamma_6 \tau_6 + (f_{60})_6 \tau + \int_0^\tau \mu_6(s_6) ds_6} d\tau$$

$$= \frac{(\rho_0)_6(0)}{\gamma_6} (e^{\gamma_6 a_6} - 1)$$

Let

$$E_6(a_6) = e^{-(f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6) d\tau_6}$$

This relationship is cardinal one. It is useful in quantum computation and other black hole dynamics problems

**Quantum space and quantum time**

(1153) **Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.**

(1154) **In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.**

(1155) **They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.**

(1156) **When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.**

(1157) **Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.**

**The Planck scale**

(1158) **Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".**

(1159) **Here ħ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.**

(1160) **The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.**

(1161) **Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)**

(1162) **Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.**

## Physical applications

**Some phenomena where quantum gravity plays a major role are the following.**

**The microscopic structure of spacetime**

(1163) **If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.**

(1164) **Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length**

(1165) **Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.**

(1166) **A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.**

**Early cosmology**

(1167) **According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.**

(1168) **However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.**

(1169) **A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.**

(1170) **Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".**

(1171) **One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.**

**Black holes**

(1172) **Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.**

(1173) **First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.**

(1174) **Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.**

(1175) **What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.**

(1176) **Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.**

**Astrophysical effects**

(1177) **Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.**

(1178) **An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.**

(1179) **The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.**

## Main approaches

(1180) **The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.**

(1181) **Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.**

(1182) **The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.**

## String theory

(1183) **String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.**

(1184) **Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.**

(1185) **String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.**

(1186) **Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.**

(1187) **However, when summed the entire series appears to be (=) divergent.**

(1188) **In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.**

(1189) **Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.**

(1190) **Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.**

(1191) **In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.**

## Loop quantum gravity and spin foams

(1192) **Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.**

(1193) **The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity.**

(1194) **Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime.**

(1195) **The granular properties of space can be (=) explicitly computed.**

(1196) **In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory.**

(1197) **The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks.**

(1198) **Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure.**

(1199) **The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit.**

(1200) **Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism.**

(1201) **This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism.**

## Noncommutative geometry

(1202) **Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime.**

(1203) **Quantum mechanics teaches us that dynamical quantities are (=) noncommutative -- meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed.**

(1204) **It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry.**

(1205) **A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces.**

(1206) **Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces.**

(1207) **This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity.**

**Others**

(1208) **Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays**

(1209) **. Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection.**

(1210) **Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible.**

(1211) **The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.**

**Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016**

(1212) **These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field.**

(1213) **The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets.**

(1214) **Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state.**

(1215) **Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands.**

(1216) **The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI**

(1217) **The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9].**

(1218) **This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive.**

(1219) **More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.**

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

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| (1220) | **Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.** |
| (1221) | **Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.** |
| (1222) | **The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.** |
| (1223) | **This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.** |
| (1224) | **The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.** |
| (1225) | **This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.** |
| (1226) | **For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.** |
| (1227) | **The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society** |

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

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| (1228) | **Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.** |
| (1229) | **We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.** |
| (1230) | **When this is true—as is often the case—the system serves as (=) a "bath" for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called thermalization.** |
| (1231) | **However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.** |
| (1232) | **Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1).** |
| (1233) | **Bringing them to thermal equilibrium requires (e) an external bath.** |

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| (1234) **But so far, there has been no rigorous proof of (e) the existence of such many-body localization. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption.** |
| (1235) **The concept of localization was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations).** |
| (1236) **The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century.** |
| (1237) **But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information.** |
| (1238) **Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment.** |
| (1239) **For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood.** |
| (1240) **It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7].** |
| (1241) **Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature.** |
| (1242) **One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.** |
| (1243) **However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].** |
| (1244) **Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.** |
| (1245) **Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.** |
| (1246) **As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked.** |
| (1247) **His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice.** |
| (1248) **An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site.** |
| (1249) **This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself).** |

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| (1250) | **This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition.** |
| (1251) | **However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime:** |
| (1252) | **These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects.** |
| (1253) | **It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies).** |
| (1254) | **This is a realistic scenario for (e) all known systems.** |
| (1255) | **On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects.** |
| (1256) | **On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other.** |
| (1257) | **No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true.** |
| (1258) | **His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found.** |
| (1259) | **Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization?** |
| (1260) | **Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects?** |
| (1261) | **Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions.** |
| (1262) | **So in those systems, the quantum phase transition does indeed exist.** |
| (1263) | **But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension.** |
| (1264) | **Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters.** |

**Focus: Biological Cells Form Electric Circuits July 1, 2016• Physics 9, 75 Cells that are electrically active and that also produce light for easy voltage monitoring could lead to new studies of heart arrhythmias and possibly bio-computing. Figure caption expand figure H. M. McNamara et al., Phys. Rev. X (2016) Cell circuit. Waves of electrical activity propagate in a loop within a collection of cells engineered with electrical and optical properties (see video below).Some sentences, pictures and videos are deleted due to constraints on space.**

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| (1265) **The human heartbeat is produced by (e) electrical pulses that propagate through (e&eb) cardiac tissue, causing (eb) rhythmic muscle contraction.** |
| (1266) **Researchers have previously engineered cells to create (eb) an artificial tissue capable of producing (eb) coordinated electrical activity, and now a team has added the ability to monitor (e&eb) their electrical state by (e) detecting fluorescent emission.** |
| (1267) **They have also fashioned the cells into (e&eb) "living circuits" that might act as (=) model systems for studying heart behavior.** |
| (1268) **Constructing electrically active tissues from their basic building blocks by (e) genetically engineering cells allows (eb) researchers to study a "clean" system similar to (e) the more complicated cardiac tissue.** |
| (1269) **These systems let them investigate, for example, the electrical signals that lead to (eb) the potentially fatal irregular heartbeat called arrhythmia.** |
| (1270) **Previously researchers have introduced (eb) ion channels—proteins that conduct ions through (e&eb) cell membranes—into human kidney cells to make (eb) electrically "excitable" tissues [1].** |
| (1271) **Adam E. Cohen of Harvard University and his co-workers decided to add more ingredients to such cells, so that the electrical activity could be (=) light-activated and easily visualized.** |
| (1272) **They incorporated four additional gene products into (e&eb) human embryonic kidney (HEK) cells.** |
| (1273) **Two were ordinary ion channels; the other two were (=) a blue-light-activated ion channel (a light-sensitive switch) and (e&eb) a protein that fluoresces with (e&eb) an intensity that depends on (=) the voltage across the membrane in which it is embedded.** |
| (1274) **Light activation of electrically excitable cells has been shown before [2]. But Cohen and colleagues have now included a convenient "readout" mechanism (the fluorescent protein) and have used patterning methods to shape (e&eb) the cell cultures into "tissue circuits."** |
| (1275) **They grew rings of cells a few millimeters in diameter that could act as (e&eb) oscillators that support (eb) circulating waves of electrical excitation.** |
| (1276) **These waves arise from (e) the collective behavior of the cells: a pulse produced by (e) one cell can trigger another from (e) its neighbors, eventually producing (eb) a coherent wave of excitation traveling through (e&eb) the tissue.** |
| (1277) **It's this coordination, in a repeating cycle of excitation and quiescence, which gives rise to (eb) rhythmic activity in cardiac tissue.** |

**H. M. McNamara et al., Phys. Rev. X (2016)**

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| (1278) **A pattern of blue light (not visible) illuminates these cells, defining (eb) a rectangular loop around which electrical waves can propagate.** |
| (1279) **At about 0:07 the rectangle begins to shrink in (eb) the vertical direction, causing t(eb) he cycle time to gradually decrease. (In other experiments, the cells were grown in (eb) fixed, ring-shaped patterns.)** |
| (1280) **The ring of engineered HEK cells acted like an electronic "ring oscillator," a device with a characteristic frequency to (e&eb) its circulating electrical signal.** |
| (1281) **The electrical waves in a ring oscillator can propagate (e&eb) in either direction around the ring, enabling (eb) it to serve as a binary memory unit.** |

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| (1282) **Cohen and colleagues found that in their HEK tissues, these circulating waves remained (eb) stable for more than two hours, with (e&eb) the pulses completing a few circuits every second.** |
| (1283) **The electrical signals were easily visualized as (=) waves of fluorescent emission.** |
| (1284) **The researchers say that these tissues could be used to (e) understand and predict the behavior of real cardiac cells, especially arrhythmias.** |
| (1285) **"In real heart, if you pace the tissue fast enough, *(eb) it starts skipping beats and eventually becomes (=) arrhythmic," says Cohen.** |
| (1286) **This state is a well-known precursor to (e) sudden cardiac death and is associated with (e&eb) the breakup of the usual regular waves of electrical excitation.** |
| (1287) **Some drugs also have the dangerous side effect of inducing (eb) arrhythmia, and Cohen says that the new system might offer a quick and convenient way of screening for such hazards. The work "is really very elegant and deserves attention," says physiologist Gil Bub of Oxford University in the UK. "It opens new research questions for people like me who are interested in wave dynamics in excitable systems."** |
| (1288) **There might also be practical uses for information processing in living tissues. "By combining patterns of cell growth and (e&eb) patterns of ion channel expression, we think we can build diodes, amplifiers, and other simple logic elements," says Cohen.** |
| (1289) **Andrew Adamatzky of the University of the West of England in Bristol, a specialist in "unconventional computing," foresees circuits that perform computations through (e&eb) the interactions of electrical waves. "This research might lead to fabrication of implantable bio-computers," he says. This research is published in Physical Review X. –Philip Ball Philip Ball is a freelance science writer in London and author of Invisible: The Dangerous Allure of the Unseen (University of Chicago Press, 2014).** |

**Optically Controlled Oscillators in an Engineered Bioelectric Tissue Harold M. McNamara, Hongkang Zhang, Christopher A. Werley, and Adam E. Cohen Phys. Rev. X 6, 031001 – Published 1 July 2016**

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| (1290) **Complex electrical dynamics in excitable tissues occur throughout (e&eb) biology, but the roles of individual ion channels can be difficult to determine due to (e) the complex nonlinear interactions in (eb) native tissue.** |
| (1291) **Here, authors ask whether they can engineer a tissue capable of (e) basic information storage and processing, where (e) all functional components are known and well understood.** |
| (1292) **Authors develop a cell line with (e&eb) four transgenic components: two to enable (eb) collective propagation of electrical waves and (e&eb) two to enable optical perturbation and (e&eb) optical readout of membrane potential.** |
| (1293) **They pattern the cell growth to define (eb) simple cellular ring oscillators that run stably for >2 h (∼104 cycles) and that can store data encoded in (eb) the direction of electrical circulation.** |
| (1294) **Using (e) patterned optogenetic stimulation authors probe the biophysical attributes of (e) this synthetic excitable tissue in detail, including (e) dispersion relations, (e&eb) curvature-dependent wave front propagation, (e&eb) electrotonic coupling, and (e&eb) boundary effects.** |
| (1295) **Authors then apply the biophysical characterization to (e) develop an optically reconfigurable bioelectric oscillator.** |
| (1296) **These results demonstrate the feasibility of engineering bioelectric tissues capable of (e&eb) complex information processing with optical input and output. Received 12 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevX.6.031001 This article is available under the terms of the** |

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**Scientists detect radio emission from a nearby brown dwarf July 6, 2016 by Tomasz Nowakowski report**

(1297) **K2 and Spitzer images of W0607+24: The position of W0607+24 is marked by (e) a red circle. SDSS J060738.06+242939.5, marked by (e) an orange circle, are the brightest star in the K2 image.**

(1298) **The three faint (r = 18.3 − 19.6) stars marked by cyan circles, contribute significant flux to (e) the K2 measurement but are (-) negligible at Spitzer's longer wavelength and (e &êb) higher resolution. Other stars with r < 20 in SDSS are marked with blue circles. Credit: Gizis et al., 2016.**

(1299) **(Phys.org)—Located some 23 light years away, the brown dwarf designated WISEP J060738.65+242953.4 or W0607+24 for short, turns out to be (=) a source of radio emission.**

(1300) **According to a research paper published July 4, this substellar object showcases (eb) quiescent radio emission, making it one of (e) the most radio-faint, ultra-cool dwarfs yet detected.**

(1301) **Brown dwarfs like W0607+24 are objects that are too large to be called planets and too small to be stars. With a mass below that necessary to maintain hydrogen-burning nuclear fusion reactions, they are much cooler and dimmer than (eb) main sequence stars.**

(1302) **These objects could be also sources of (e) radio emission, but the cause of this process is still not completely understood.**

(1303) **Due to its proximity, W0607+24 is (=) an excellent target for studying the nature of (e) brown dwarfs.**

(1304) **It is the nearest known late-L dwarf in the northern hemisphere, and the third-nearest in (=) the whole sky, (eb) allowing researchers to study phenomena occurring in brown dwarfs.**

(1305) **That's why a team of astronomers, led by John Gizis of the University of (e)Delaware, Newark, carried out a series of mid-infrared, radio, and optical observations of W0607+24 to study any activity in this curious substellar object.**

(1306) **Motivated by the results from NASA's prolonged Kepler mission, named K2, which obtained (eb) long time-series photometry of W0607+24; they observed this object simultaneously with the Karl G. Jansky Very Large Array (VLA) in New Mexico and NASA's Spitzer Space Telescope.**

(1307) **The most important results were provided by VLA, as the team found that W0607+24 is (=) a source of quiescent, non-bursting, radio emission.**

(1308) **The scientists noted that the discovery of this process in this nearby object could be (=) crucial for our understanding of this process in (eb) ultracool dwarfs in general.**

(1309) **"W0607+24 is a valuable test case for (e) understanding the mechanisms of radio emission in ultracool dwarfs.**

(1310) **Some of the radio-detected objects emit (eb) rapid, intense bursts with high circular polarization, often found to occur periodically and attributed to coherent emission due to (e) the electron cyclotron maser instability," the paper reads.**

(1311) **The astronomers also suggested that the relative radio faintness of W0607+24 may be due to (e) its orientation if (e) the emission is concentrated at the poles.**

(1312) **However, follow-up observations are needed to better characterize (eb) the detected radio emission and to eliminate (e) all the uncertainties.**

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| (1313) **The team also hopes for more significant findings when (e) studying this nearby substellar object in the future.** |
| (1314) **"As the nearest known northern hemisphere late-L dwarf, this offers many possibilities for future studies. For example, it could serve as (=) an useful standard for (e) line-broadening and (e&eb) polarization studies.** |
| (1315) **If there is a planetary system aligned with (e&eb) (the primary's rotation axis, then we would expect very little radial velocity signal but a strong astrometric signal," the researchers concluded.** |
| (1316) **Moreover, the team plans to monitor (e&eb) additional late-M and L dwarfs in other K2 campaigns, expecting even more promising results.** |
| (1317) **They are more distant than W0607+24, but because they are warmer, (eb) some are brighter and higher signal-to-noise than that described in this paper.** |
| (1318) **In addition to detecting the radio emission, (eb)the scientists also estimated the age, mass and radius of W0607+24** |
| (1319) **. According to the study, the object is less than two billion years old, has (e) a maximum mass of about 0.055 solar masses and a radius approximately of 0.1 solar radii.** |
| (1320) **More information: WISEP J060738.65+242953.4: A Nearby. Pole-On L8 Brown Dwarf with Radio Emission, arXiv: 1607.00943 [astro-ph.SR] arxiv.org/abs/1607.00943** |
| (1321) **Authors present a simultaneous, multi-wavelength campaign targeting the nearby (7.2 pc) L8/L9 (optical/near-infrared) dwarf WISEP J060738.65+242953.4 in (eb) the mid-infrared, radio, and optical.** |
| (1322) **Spitzer Space Telescope observations show no (e) variability at the 0.2% level over 10 hours each in the 3.6 and 4.5 micron bands.** |
| (1323) **Kepler K2 monitoring over 36 days in Campaign 0 rules out (e) stable periodic signals in (eb) the optical with amplitudes great than 1.5% and periods between 1.5 hours and 2 days.** |
| (1324) **Non-simultaneous Gemini optical spectroscopy detects lithium, constraining (e) this L dwarf to be less than ~2 Gyr old, but no Balmer emission is observed.** |
| (1325) **The low measured projected rotation velocity (v sin i < 6 km/s) and lack of variability are (=) very unusual compared to other brown dwarfs, and authors argue that this substellar object is (=) likely viewed pole-on.** |
| (1326) **Authors detect quiescent (non-bursting) radio emission with (e&eb) the VLA.** |
| (1327) **Amongst radio detected L and T dwarfs, it has (e) the lowest observed L_nu and the lowest v sin i.** |
| (1328) **They discuss the implications of pole-on detection for (e) various proposed radio emission scenarios. © 2016 Phys.org** |

**A new look at the galaxy-shaping power of black holes July 6, 2016**

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| (1329) **It wasn't until the advent of X-ray astronomy that the full picture of (e) the galaxy-shaping power of black holes began to emerge (eb) with the ability to see plasma.** |
| (1330) **Data from a now-defunct X-ray satellite is providing (eb) new insights into the complex tug-of-war between galaxies, the hot plasma that surrounds them, and the giant black holes that lurk in their centres.** |

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| (1331) | **Launched from Japan on February 17, 2016, the Japanese space agency (JAXA) Hitomi X-ray Observatory functioned for just over a month before contact was lost and (e&eb) the craft disintegrated.** |
| (1332) | **But the data obtained during those few weeks was enough to paint (eb) a startling new picture of the dynamic forces at work within galaxies.** |
| (1333) | **New research, published in the journal Nature today, reveals data that shows just how important the giant black holes in galactic centres are to the evolution (eb) of the galaxies as a whole.** |
| (1334) | **"We think that supermassive black holes act(e&eb) like thermostats," said Brian McNamara, University Research Chair in Astrophysics at the University of Waterloo. "They regulate (e&eb) the growth of galaxies."** |

**Champagne bubbles of plasma**

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| (1335) | **During its brief life, the Hitomi satellite collected X-ray data from (e) the core of the Perseus cluster, an enormous gravitationally-bound grouping of (e) hundreds of galaxies.** |
| (1336) | **Located some 240 million light years from earth, the Perseus cluster is one of (e) the largest known structures in the universe.** |
| (1337) | **The cluster includes not only the ordinary matter that makes up the galaxies, but an (e) "atmosphere" of hot plasma with a temperature of tens of millions of degrees, as well as a halo of (e) invisible dark matter.** |
| (1338) | **Earlier studies, going back to the 1960s, have shown that each of the galaxies in the cluster - and indeed most galaxies - likely contains (e) a supermassive black hole in its centre, an object 100 million to more than ten billion times as massive as our sun** |

**New look at the galaxy-shaping power of black holes**

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| (1339) | **Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving.** |
| (1340) | **"These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science.** |
| (1341) | **"Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy."** |
| (1342) | **This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return.** |
| (1343) | **The remaining matter gets absorbed into the black hole, adding to (e) its mass.** |
| (1344) | **The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne.** |
| (1345) | **The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution.** |
| (1346) | **Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The** |

**answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form.**

(1347) **"Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara.**

(1348) **"It creates (eb) these bubbles, and the bubbles keep the plasma hot.**

(1349) **That's what prevents galaxies from (e) becoming even bigger than they are now."**

(1350) **Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.**

(1351) **In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.**

(1352) **Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.**

(1353) **"It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.**

(1354) **"It's like nature's thermostat, which keeps these galaxies from (e) growing.**

(1355) **If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."**

**A new look at the galaxy-shaping power of black holes**

(1356) **An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the**

(1357) **McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.**

(1358) **Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.**

**Doomed satellite missions**

(1359) **Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.**

(1360) **At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.**

(1361) **Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a**

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| **leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.** |
| (1362) **The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.** |
| (1363) **These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.** |
| (1364) **"But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people** |
| (1365) **. It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."** |
| (1366) **The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.** |
| (1367) **"This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website** |

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

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| (1368) **Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.** |
| (1369) **They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.** |
| (1370) **However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.** |
| (1371) **Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.** |
| (1372) **X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well.** |
| (1373) **The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma.** |
| (1374) **These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6.** |
| (1375) **Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus.** |

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| (1376) | **A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate.** |
| (1377) | **Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure.** |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (1378) | **This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy.** |
| (1379) | **Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole.** |
| (1380) | **But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy.** |
| (1381) | **A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole.** |
| (1382) | **It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth.** |
| (1383) | **"There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA).** |
| (1384) | **This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies.** |
| (1385) | **However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated.** |
| (1386) | **In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1.** |
| (1387) | **We know dark matter exists only from (e) its gravitational effects.** |
| (1388) | **It holds together galaxies and (e&eb) galaxy clusters.** |
| (1389) | **Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years.** |
| (1390) | **To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies.** |
| (1391) | **They used star motions as tracers to (e) weigh the galaxies' central black holes.** |
| (1392) | **X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto.** |
| (1393) | **They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone.** |
| (1394) | **This connection is likely to be related to (e&eb) how elliptical galaxies grow.** |

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| (1395) **An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together.** |
| (1396) **Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole.** |
| (1397) **"In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp** |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (1398) **The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the …** |
| (1399) **Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness?** |
| (1400) **Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness.** |
| (1401) **"Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain."** |
| (1402) **Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness.** |
| (1403) **"With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication."** |
| (1404) **Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain** |

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| **changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories.** |
| (1405) **The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen.** |
| (1406) **In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer.** |
| (1407) **Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants.** |
| (1408) **These were treated as (=) "aware" and "unaware" trials, respectively.** |
| (1409) **Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it.** |
| (1410) **The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another.** |
| (1411) **Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion.** |
| (1412) **No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness.** |
| (1413) **"We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said.** |
| (1414) **"Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness."** |
| (1415) **The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner.** |
| (1416) **Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain.** |
| (1417) **"We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said.** |
| (1418) **"This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients** |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

(1419) **Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter.**

(1420) **As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined.**

(1421) **Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions.**

(1422) **"Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.**

(1423) **The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.**

(1424) **However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.**

(1425) **The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.**

(1426) **Noda and his colleagues instead studied the change in emissions over (e&eb) time.**

(1427) **By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.**

(1428) **Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).**

(1429) **The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.**

(1430) **Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.**

(1431) **"The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp**

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(1432) **The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking "palm tree" organization.**

(1433) **The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society**

(1434) **(Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible?**

(1435) **With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem.**

(1436) **This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other.**

(1437) **According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible.**

(1438) **This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges.**

(1439) **In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century.**

(1440) **Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching.**

(1441) **In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom.**

(1442) **Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth.**

**Spheres on trees**

(1443) **Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration.**

(1444) **They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement.**

(1445) **Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum.**

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| (1446) **By visualizing the problem in this way, the researchers noticed that these particular graphs don't (e) have lots of branches extending from lots of (e) other branches.** |
| (1447) **Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure.** |
| (1448) **The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network.** |
| (1449) **As a result, it provides (eb) an important clue for finding the global minimum.** |
| (1450) **It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum.** |
| (1451) **This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large.** |

**Thompson's problem 2**

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| (1452) **By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society** |
| (1453) **The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima.** |
| (1454) **The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node.** |

**Implications for Thomson's and Smale's problems**

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| (1455) **Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before.** |
| (1456) **"Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around."** |
| (1457) **As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem.** |
| (1458) **"I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view."** |
| (1459) **In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges.** |
| (1460) **From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear.** |

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| (1461) | **"While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK.** |
| (1462) | **"For larger numbers, there will likely be (=) multiple funnels.** |
| (1463) | **Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum."** |
| (1464) | **Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.** |
| (1465) | **"The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.** |
| (1466) | **"Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org** |

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**Introduction:**

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price, University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss.**

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| **(193)** | **Black Holes are regions of (e) space in which gravitational fields are so strong that no particle or signal can escape (eb) the pull of gravity.** |
| **(194)** | **The boundary of this no-escape region is (=) called the event horizon, since distant observers outside the black hole cannot (e) see (cannot get light from) events inside.** |
| **(195)** | **Although the fundamental possibility of such an object exists within (eb) Newton's classical theory of gravitation, Einstein's theory of gravity makes (eb)black holes inevitable under some circumstances.** |
| **(196)** | **Prior to the early 1960s, black holes seemed to be (=) only an interesting theoretical concept with no (e) astrophysical plausibility, but with (e&eb) the discovery of quasars in 1963 it became clear that very exotic astrophysical objects could exist.** |
| **(197)** | **Nowadays it is taken for granted that black holes do exist in (eb) at least two different forms.** |
| **(198)** | **Stellar mass black holes are (=) the endpoint of the death of some stars, and supermassive black holes are the result of (e) coalescences in the centers of most galaxies, including our own.** |
| **(199)** | **No signal can propagate from (e) inside a black hole, but the gravitational influence of a black hole is (=) always present.** |

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| **(200) This influence does not (e) propagate out of the hole; it is permanently present outside, and depends only on (=) the total amount of mass, angular momentum, and electric charge that have gone into forming (eb) the hole.** |
| **(201) Black holes can be detected through the influence of this (e&eb) strong gravity on the surroundings just outside the hole.** |
| **(202) In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them.** |
| **(203) Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors.** |

**Classical vs. relativistic black holes**

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| **(204) Something like a black hole exists within (eb) Newton's classical theory of gravity.** |
| **(205) In that theory, an energy argument tells us that there is (=) an escape velocity $v_{esc}=\sqrt{2GM/R}$ from (e) the surface of any spherical object of (e) mass M and radius R** |
| **(206) If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity.** |
| **(207) Thus the condition for such an "unseeable" object is $R<2GM/c^2$** |
| **(208) In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape.** |
| **(209) This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity.** |
| **(210) To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it.** |
| **(211) The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c .** |
| **(212) This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed.** |
| **(213) An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity.** |
| **(214) Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable.** |

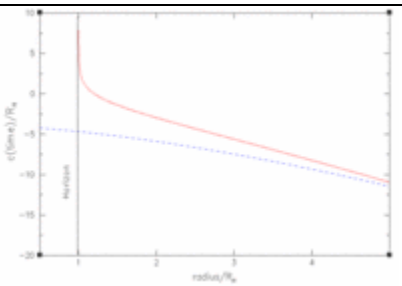
**Stationary black holes**

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| **(215) In the Newtonian theory, gravity is described by (e) the potential $\Phi$.** |
| **(216) Inside a spherical object the form of (e) $\Phi(r)$ depends on (=) the interior structure, but in the vacuum outside (e) matter the potential $-GM/r$ depends only on (=) the interior mass.** |
| **(217) Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object.** |

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| **(218) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward.** |
| **(219) This horizon radius RH=2GM/c2 is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory.** |
| **(220) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution.** |
| **(221) Radius RHhere actually means that the area of the event horizon is (=) 4πR2H.** |
| **(222) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution.** |
| **(223) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity.** |
| **(224) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own.** |
| **(225) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole.** |
| **(226) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density.** |
| **(227) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed.** |
| **(228) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole.** |
| **(229) In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter.** |
| **(230) Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum.** |
| **(231) In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior.** |
| **(232) The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object.** |
| **(233) Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum.** |
| **(234) These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.]** |
| **(235) All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time.** |
| **(236) But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?"** |
| **(237) The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point.** |

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| **(238) One of these coordinates is (=) called the "coordinate time."** |
| **(239) Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate.** |



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| **(240) Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole** |
| **(241) At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole.** |
| **(242) Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations.** |
| **(243) For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree** |
| **(244) An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events.** |
| **(245) The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving.** |
| **(246) Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle.** |
| **(247) The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return.** |
| **(248) But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon.** |
| **(249) Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future.** |
| **(250) The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?"** |

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| **Black hole parameters** |

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| **(251) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin).** |
| **(252) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun.** |
| **(253) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits.** |
| **(254) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form.** |
| **(255) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds.** |
| **(256) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2.** |
| **(257) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed.** |
| **(258) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer.** |
| **(259) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2)** |
| **(260) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials.** |
| **(261) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable** |
| **(262) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite.** |
| **(263) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole.** |
| **(264) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars.** |
| **(265) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity.** |
| **(266) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity.** |
| **(267) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite.** |
| **(268) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity.** |
| **(269) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum.** |

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| **(270) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship.** |
| **(271) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= (4πG2/c4)[(M+M2−a2−−−−−−−√)2+a2].** |
| **(272) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored.** |

**Dynamical black holes**

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| **(273) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape.** |
| **(274) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, a horizon can be born.** |
| **(275) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid.** |
| **(276) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward.** |
| **(277) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations.** |
| **(278) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes.** |
| **(279) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3.** |
| **(280) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/** |
| **(281) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers.** |
| **(282) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission.** |
| **(283) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.** |
| **(284) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.** |
| **(285) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.** |
| **(286) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.** |

**(287) The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature 6×10−8 (Mo/M) K.**

**(288) Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.**

Then $[Y_6]_6[\gamma_6]_6, a_6) = E_6(a_6)e^{-\gamma_6 a_6}$. by a simple calculation,

$$\int_0^\infty \beta_6(a_6)Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma}\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}\ da\right);$$

$$\int_0^\infty Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma_6}\left(\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}\ da_6\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_6(0)}{\gamma_6}[f_{60}^6]\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}da_6\ (\int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6)$$

$$+\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right).$$

$$\left(1 + \frac{(\rho_0)_6(0)}{\gamma_6}(f_{60}^{'})_6\left(\int_0^\infty E_6(a_6)da_6 - \int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}da_6\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right)(1 + \frac{(f_{60}^{'})_6}{\gamma_6}(\rho_0)_6(0)\int_0^\infty E_6(a_6)da_6) = 0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1 - \int_0^\infty \beta_6(a_6) Y_6(\gamma_6 a_6) da_6 = 1 - \int_0^\infty \beta_6(a_6) E_6(a_6) e^{-\gamma_6 a_6} da_6 \neq 0,$$

Since $\int_0^\infty \beta_6(a_6) Y_6(\gamma_6 a_6) \int_0^a \alpha_6(\tau_6) Y^{-6}(\gamma_6, \tau_6) d\tau_6 da_6 > 0$. Therefore, the characteristic solution is

$$\gamma_6 = -(f_{60}^{'})_6 (\rho_0)_6 \int_0^\infty E_6(a_6) da_6.$$

Since $f_6'(P_0)_6 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_6), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE OF DEBIT TYPES AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_6), 0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_6 + (\mu_6(a_6) + [f_6]_6(P_6^*) + [f_2]_6(Q_6^*))P_6^* = 0, \quad (26)$$

$$p_6^*(0) = h_6(Q_6^*) \int_0^\infty \beta_6(a_6) p_6^*(a_J) p_6^*(a_6) da_6 \quad (27)$$

$g_6(P_6^*) - b_6 = 0,$ (28)

Where $P_6^*$=$\int_0^\infty p_6^*(a_6)\, da_6$

Equation (28) yields $P^* = g^{-6}(b)$,

From (26),

$p_6^*(a) = p_6^*(0)e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6}$

Substituting into (27), we need to solve

$$\left(1 - h_6(Q_6^*)\int_0^\infty \beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6} da_6\right) p_6^*(0) = 0$$

For $Q_6^*$

Fix $p_6^*(0)$ and $P_6^*$ Let

$$H_6(z_6) = 1 - h_6(z_6)\int_0^\infty \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6} da_{,6}$$

And assume that

$$[k_3]_6 \int_0^\infty \beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(g_6^{-6}(b_6))a_6 - [k_2]_6 a_6}\, da_6 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_6(z_6)] > 0.$$

Clearly,

$$\frac{dH_6}{dz_6} = h_6(z)([f_2^{'}]_6(z_6))\int_0^{\infty} a_6\beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - f_6(P_6^*)a_6 - f_2(z_6)a_6}\, da_6$$

$-h^{'}(z_6)\int_0^{\infty}\beta_6(a_6)\, e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6}\quad da_6$

By the assumptions on $\left[f_2^6\right]_6 and\ [h^6]^6 dH_6/dz > 0$. Hence, if

$$H_6(0) = 1 - \int_0^{\infty}\beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6}da_6 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_6(0) < 0$ follows if $g_6^{-}(b_6) < [f_6^{-6}]_6(r_6)$, where $r_6$ is the solution of (12).

Indeed, since

$$R_6(s_6) = \int_0^{\infty}\beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 e^{-y}6 a}da_6$$

Is a decreasing function of $s_6$ when $P_6^* = [g^{-6}]_6(b_6) < [f_6^{-6}]_6(r_6)$, that is $[f_6]_6(P_6^*) < r$, $R\left(f_6(P^*)\right) > 1$which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_6(g^{-6}(b))) > 6, i.e.,$

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6\left(g_6^{-6}(b)\right)a_6} > 1. \quad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_6^*(a_6), Q_6^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_6^*(a_6), Q_6^*)$ is more complicated. We again consider perturbations $m_6(t_6 a_6) = p_6(t_6, a_6) - p_6^*(a_6)$ and $W_6(t) = Q_6(t) - Q^*$.Then

$$M_6(t) \equiv \int_0^\infty m_6(t, a_6) da_6 = P_6(t) - P_6^*,$$

$$[m_6]_6 + [m_a]_6 = [p_6]_6 + [p_a]_6 - \left[p^*\right]_6(a_6) \approx -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6$$

$$-[f_{6*}^6]_6 p_6^*(a_6)M - [f_{2*}^6]_6 p_6^*(\mathrm{a})W_6,$$

$$m_6(t,0)p_6(t,0) - p_6^*(0) \approx h_* \int_0^\infty \beta_6(a_6) m_6(t, a_6)\, da + \frac{[h_*^6]_6}{[h_*]_6}\, p_6^*(0) W_6$$

$$\frac{dW_6}{dt} = \frac{dQ_6}{dt} \approx [g_*^6]_6 Q_6^* M_{,6}$$

Where $[f_{6*}]_6$ denoting $[f_6(P^*)]_6$, $[f_{6*}]_6$ denoting $[f_2(P^*)]_6$, $\left[f_{6*}^{'}\right]_6$ denoting $\left[f^{'}{}_6(P^*)\right]_6 \left[f_{2*}^{'}\right]_6$ denoting $[f_2]_6$ $(Q_6^*,$ $h_{*}^{'}$ denoting $h_{6}^{'}$ $(Q_6^*)$, and $\left[g_*^{'}\right]_6$

Denoting $\left[g^{'}\right]_6(P_6^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_6]_6 + [m_a]_6 = -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6 - [f'_{6*}]_6[p_6^*](a_6)M_6 - [f'_{2*}]_6 p_6^*(a_6)W_6$$

$$m_6(t,0) = [h_*]_6 \int_0^\infty \beta_6(a_6)m_6(t,a_6)da_6 + \frac{[h_*^6]_6}{[h^*]_6}[p^*]_6(0)\,W_{,6} \qquad (30)$$

$$\frac{dW_6}{dt} = [g_*^6]_6[Q_6^*]M_6$$

Integrating $[m_6]_6(t,a_6)$ with respect to a from 0 to ∞, we have

$$\frac{dM_6(t)}{dt} = \int_0^\infty [mt]_6(t,a_6)da_6$$

$$= -\int_0^\infty [m_a]_6\,(t,a_6)da_6 - \int_0^\infty \mu(a_6)m(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + [f'_{6*}]_6 \int_0^\infty p_6^*(a_6)da_6\right)M_6 - [f_{2*}^6]_6 \int_0^\infty p_6^*(a_6)da_6 W_6$$

$$= m_6(t,0) - \int_0^\infty \mu_6(a_6)\, m_6(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + [f'_{6*}]_6 P_6^*\right)M_6 - [f'_{2*}]_6 P_6^* W_{,6}$$

From (30), we arrive at the equations

$$\frac{dW_6(t)}{dt} = \int_0^\infty ([h_*]_6\beta_6(a_6) - \mu_6(a_6))m_6(t,a_6)da_6$$

$$-\left([f_{6*}] + [f_{2*}]_6 + [f'_{6*}]_6 P_6^*\right)M_6 + \left(\frac{[h'_*]_6}{[h_*]^6}\,P_6^*(0) - f'_{2*}P_6^*\right)W_6,$$

(36)

$$\frac{dW_6(t)}{dt}\left[g^{'}{}_{*}\right]_6 Q_6^{*} M_{.6}$$

By the mean value theorem,

$$\int_0^{\infty}([h_{*}]_6\beta_6(a_6)-\mu_6(a_6))m_6(t,a_6)da_6=([h_{*}]_6\beta_6(\theta_6)-\mu_6(\theta_6))M_{,6}$$

Where $0<\ \theta_6=\theta_6(t)<\infty$.

And if $a_6^2\,[a_6^2]_6<4\,[a_2]_6<4\,a_2 e$,

$\mathrm{Re}\,\lambda_6\left[\lambda_6\right]_6=\mathrm{Re}\lambda_2\left[\lambda_2\right]_6=-\frac{6}{2}a_6[a2\ ]_6<0.$

Let $\Delta=\max_{2<u<\infty}\left|h_{*}\,[h_{*}]_6\beta_6(a_6)-\mu_6(a_6)\right|$. Then$|n_6(t)|$ 2Δ∀t 0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^{*}(a),\boldsymbol{Q}^{*})$ is asymptotically stable if either $a_6^2<4\,a_2 e$ and $[a_6]>2\Delta$,or $[a_6^2]$ 4$[a_2]$e and $(a_6-\sqrt{(a_6^2)_6-4(a_2)_6 e})>2\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a)=\mu_0+\left(\frac{a}{a+6}\right)$ and $\beta_6(a_6)=ae^{-ma}$ Then $\Delta=[\mu_0]+1$ Simple calculation leads to

$$\int_0^{\infty}\beta_6(a_6)e^{-\int_0^a\mu_6(\tau_6)+d\tau_6}da=\frac{6}{(m+\mu_0+1)^2}+\frac{2}{(m+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_6 + [\mu_0]$ 0.5. Assume$[f_6](x_6)= K_6x_6$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_6(x_6)= [0x_6]/(1 + x_6)$; the functional responses $[[f_2](x) = [0]\, x_6/(1 + x_6)$and $h_6(x_6)$=(1+$\omega_6 x_6$)/(1 + $x_6$)(See Holling [7] and Rashevsky [13])

Equation (64),now has the form of

$$\frac{m_6 + [\mu_0] + [\mathrm{K}][P_0])}{[m_6]_{\,6} + (\mu_0)_{\,6} + [\mathrm{K}][P_0])\,{}^{3}_{6}\ [\mathrm{K}]_{\,6}[P_0]\,{}^{3}_{6}} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(6 + P_0) < b_6$i e $[P_0] < b/(\delta_6 - b_6)$then the equilibrium$(P_0)$ $[P_0]_{\,6}(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_6$

$[f_{6*}] = K_6\ \mathrm{P}^{*}\ [f_{2*}] = \frac{\sigma \mathrm{Q}_6^*}{6+\mathrm{Q}_6^*}$ $[h_*] = \frac{6+\omega_6 \mathrm{Q}_6^*}{6+\mathrm{Q}_6^*}$

$[f_{6*}] = K_6\ \left[\mathrm{f}\,{}^{6}_{2*}\right] = \frac{\sigma_6}{(6+\mathrm{Q}_6^*)^2}[\ \mathrm{g}\,{}^{6}_{*}] = \frac{\sigma_6}{(6+\mathrm{P}_6^*)^2}$ $\left[\mathrm{h}\,{}^{6}_{*}\right] = \frac{\omega_6 - 6}{(6+\mathrm{Q}_6^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0$, K, b, δ, and $\omega$ need to satisfy

$$\frac{\omega_6}{(m + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 + b_6} + \sigma)^{\,2}} + \frac{2\omega_6}{([m_6]_{\,6} + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 + b_6} + \sigma_6)^{\,3}} < 1$$

And

$$\frac{1}{(m_6 + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 - b_6})^{\,2}} + \frac{2\omega_6}{(m_6 + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 - b_6})^{\,3}} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_6 + [\mu_0]_6 + 1 + \frac{K_6 b_6}{\delta_6 - b_6} + (m_6 + [\mu_0]_6 + \frac{K_6 b_6}{\delta_6 - b_6} + \sigma_6) Q_6^*)^3$$

$$= m_6 + [\mu_0]_6 + 3 + \frac{K_6 b_6}{\delta_6 - b_6}(1 + Q^*)(1 + \omega_6 Q_6^*$$

Now

$$\left[a_6^2\right] - 4\left[a_2\right]e = \left(\frac{2K_6 b_6}{\delta_6 - b_6}\right)^2 + \frac{4K_6 b_6 \sigma_6 Q_6^*}{(\delta_6 - b_6)(1 + Q_6^*)} + \frac{\sigma_6^2 Q_6^{*2}}{(1 + Q_6^*)^2}$$

$$\frac{4\, b_6 \sigma_6 (\delta_6 - b) Q^*}{\delta_6 (1 + Q_6^*)^2} - \frac{4 (\delta_6 - b_6)^2 (1 - \omega_6) Q_6^*}{\delta_6 (1 + \omega_6 Q_6^*) 1 + Q_6^*)} P * (0)$$

And

$$\Delta = \lim_{0 < a < \infty} |[h_*]_6 \, ae_6^{mu} \; [\mu_0]_6 - \frac{a_6}{a_6 + 6}$$

Hence if $[a_6^2 > 4(a_2)_6 e$ and

$$[a_6] = \frac{2K_6 b_6}{\delta_6 - b_6} + \frac{\sigma_6 Q_6^*}{6 + Q_6^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_6]_6 = [f_{6*}]_6 + [f_{2*}]_6 + [f_{6*}]\,P^*_6$$

$$[a_2]_6 = \left[f^6_{2*}\right]P^*_6 - \frac{h^6_*}{h_*}P^*(0)$$

$$e=[g^6_*]_6\,Q^*_6$$

$$\eta_6(t) = \left[h_*\right]_6\beta_6\left(\theta_6(t)\right) - \mu_6\left(\theta(t)\right),$$

The system(36)can be written as the following system:

$$[M(t)]\left[a_6\right]M_6 - [a_2]_6 W_6 - [n(t)]M_6$$

$$w_6(t) = eM_6$$

Write

$$A=\begin{pmatrix} -[a_6]_6 - [a_2]_6 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_6 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_6 \\ W_6 \end{pmatrix}$$

The solution of (32)is

$$X_6(t) = e^{A,t}\,[X]_6(0) + \int_0 e^{At}$$

Other problems to which the model is applied are: To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case β→∞ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness condition of (e) entropy. It is shown that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology

(gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure authors have discovered. It is shown that Legendre invariant thermodynamic scalar curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws, leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions,**

**geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged AdS black holes by considering the case of the RN-AdS black holes surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase transition. For the case $\omega_q=-2/3$ authors derive analytic expressions of (e) critical physical quantities, while (e) for cases $\omega_q\neq-2/3$ appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark

energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=)

chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross

section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions, while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center

collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma? Paul M. Chesler STOP Paul M. Chesler, Krishna Rajagopal** calculate how the energy and (e&eb) the opening angle of jets in N=4 SYM theory evolve as they propagate through (e&eb) the strongly coupled plasma of that theory. **Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb)

heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version) **Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb) strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion

collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar (=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI: 10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills

(SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the

decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed

black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to

equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a

nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about

the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic

map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstates thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11 August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical

advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.**

Characterizing the behaviour of strongly coupled quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at $\lambda=\lambda c$, the correlation length diverges as $\xi\sim|\lambda-\lambda c|-\nu$, while (e) the gap between the ground state and first excited state

vanishes as (e) ΔE∼ξ−z∼|λ−λc|zν. Exactly at the critical point λ=λc, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as ω∼|k−k0|z. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004 STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting

qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element. The basic architecture presented here can be expanded, offering (eb) the possibility for the coherent interaction of (e&eb) a large number of superconducting qubits. **Nature 449, 438-442 (27 September 2007) | doi:10.1038/nature06124; Received 18 April 2007; Accepted 25 July 2007 Coherent quantum state storage and transfer between two phase qubits via a resonant cavity Mika A. Sillanpää1, Jae I. Park1 & Raymond W. Simmonds1 STOP. Chui-Ping Yang, Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e)

long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007**

# PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION ATTRITION MODELS

## MOODULE NUMBER SEVEN

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Systems to which have applied the model are as follows: Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures have been proposed. An attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in**

**(eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP. In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011 http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 Qubit transfer**

between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys. Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer. Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP. In recent years S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P have investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys.

**Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP. Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2 [quant-ph] for this version) Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP. Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue from over (e&eb) a thousand plant and animal samples. NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP. Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this behavior, enhancing and accelerating (eb, eb+) the growth of correlations. Michael Kastner gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. Michael Kastner tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP. A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E.**

**Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91, 075435 – Published 26 February 2015 STOP. Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1 proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point'' with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. Membranes at quantum criticality Petr Hořava1, 2 Published 4**

**March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that (eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035 Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013 STOP Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study**

**the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen STOP Authors study the dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing**

**seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) Dynamic Service Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by H∗(X,(Λ∗TX)⊗g), where X is the target hyperk\"ahler manifold. We further give a mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the**

**low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition. Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb) the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537 © 2016 American Physical Society Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant (Δα/α) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and Δα/α to (e&eb) this conclusion.**

**However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts z≲3, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the Λ cold dark matter model value (jΛCDM=1); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in α(z), linked to (e&eb) a linear coupling (ζ) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order |ζ|∼10−4, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in Δα/α (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems, authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them**

**are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP Authors measure the star formation properties of (e) two large samples of galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the star formation rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a function of environment. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]. Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size √ n × √ n? The problem is that √n time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time O (√n) for d ≥ 3, or O (√nlog5/2n) for d = 2. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ∗This work was mostly done while the author was a PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship and by CIAR. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In**

**d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e) sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP The relationship between the foundations of mathematics and (e&eb) physics is a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb)**

**flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph] (or arXiv: 1503.05600v1 [math-ph] for this version) Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative**

**real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to the fine structure constant. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor goes to 0 as the (e&eb) time goes to 0, the big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph]**

**(or arXiv: 1306.4613v1 [quant-ph] for this version) Effects of gauge theory based number scaling on geometry Paul Benioff STOP In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x, outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph] (or arXiv: 1211.3381v2 [quant-ph] for this version) Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value, \int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances, the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an**

**observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include (e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP This work is based on (e) the field of reference frames based on (e) quantum representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul Benioff STOP A quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of**

**these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference: Journal of Physics: Conference Series 70 Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows: System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The**

**Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol**

## 23. MODEL FORMULATION

Let $p_7(t, a_7)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(g_7)Q_7$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_7(0) = 0, \quad g_7(P_7) > 0, \ P_7 \ 0, \qquad \lim_{P_7 \to \infty} g_7(P_7) = (K_7)_7 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_7 + (\rho_a)_7 + D_7(a_7, P_7, Q_7)\rho_7 + 0, \qquad 0 < a_7 \ \infty,$$

$$\rho_7(t,0) = \int_0^\infty n_7(a_7, P_7, Q_7)\rho_7(t_7, a_7)da_7$$

Where $D_7$the mortality or debit type function is function and $n_7$ is the fertility function.

Put $_7$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_7}{dt}\right]_7 = x_7 \left[\frac{dQ_7}{dt}\right]_7 - x_7 \left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_7}{dt}\right]_7 = x_3\left[\frac{dQ_3}{dt}\right]_7 - x_4\left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_7}{dt}\right]_7 - x_6\left[\frac{dQ_3}{dt}\right]_7$$

Where $x_7, x_7, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_7$ and $g_7(P_7)Q_7$

Now when you add the term $(b_7Q_7)$ with $g_7(P_7)(Q_7)$ we are infect see is that $(b_7Q_7)$ is credited incrementally part of $g_7(P_7)(Q_7)$ or instant nearly, when subtracted from $\frac{dQ_7}{dt}$ , it means that $g_7(P_7)(Q_7)$ is slowly eaten by $\frac{dQ_7}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

Following equation throws light on important issues in the following cases

**Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA**

(1467) (Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor.

(1468) In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface.

(1469) As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone.

(1470) Such planets are all the "right" distance from (e) their star, and hopefully, also have water.

(1471) But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone.

(1472) For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement.

(1473) But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway.

(1474) But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain.

(1475) It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life.

(1476) Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface.

(1477) The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113

**Abstract**

(1478) The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question.

(1479) Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry.

(1480) This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry.

(1481) Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments.

(1482) Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy.

(1483) The electronic and (e&eb) structural variability among them is extraordinary.

(1484) The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations.

(1485) The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface.

(1486) The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers.

(1487) Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed

mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp

**Water waves travel back in time to retrace their ripples By Jacob Aron (New Scientist)**

(1488) Nudge your bath tub just right and (eb) you might end up reversing time.

(1489) Don't worry – you won't end up wet, naked and without a towel in the Cretaceous period. Emmanuel Fort at the Langevin Institute in Paris, France, and his colleagues have developed what they call a "time mirror", which can make water waves rewind into (e&eb) their past. Humans and rubber duckies are (=) affected.

(1490) To understand how this works, imagine dropping a stone in a pond and watching waves ripple out from a central point in an ever-expanding circle. The equations governing the behaviour of waves say that it should also be possible for a ring of ripples to refocus back into a single point, like a reversed movie of the dropped stone. But we never actually see this happening, because it is vanishingly unlikely that an outside-in ring would form spontaneously. Put it in reverse

(1491) Fort and others have previously shown that it is possible to engineer the situation so the waves turn back on themselves by kitting out the edge of your pond with sensors and wave generators, effectively recording and then playing back the waves in reverse. The process has been done with water, sound and even microwaves, and is used to produce clearer ultrasound images for medical uses.

(1492) But now the team has discovered that simply dropping the tub downwards is enough to send water waves scurrying in reverse, with no need for complex sensors. That's because the velocity of water waves depends on acceleration due to gravity. "If you do a sudden jolt, you change the effective gravity, and the velocity of the wave changes," says Fort.

(1493) The disruption essentially resets the starting point of the wave. Once normal gravity is restored, it starts to travel both backwards and forwards. "Suddenly, you freeze the wave. When you release it again, it doesn't remember that it was going one way," says Fort.

(1494) Ripples in the water returned exactly to their start point - the shape of the Eiffel Tower

(1495) Ripples in the water returned exactly to their start point – the shape of the Eiffel Tower

**Déjà vu**

(1496) To show that the backward-travelling waves converge exactly on their starting point, just as if you were rewinding time, the team conducted a series of experiments. In some, they simply created a circular ripple of waves by dipping a probe in the water, similar to dropping a stone in the pond. In others, they used wave emitters in the shape of France, the Eiffel Tower and a smiley face to create complex wave-fronts. In each case, the reversed waves converged to recreate their original shape.

(1497) The experiments aren't truly reversing time, because the waves continue to travel out again once they have converged at the central point, points out Fort. "A real time-reversal would have a sink that absorbs all the waves," he says.

(1498) Now the team hopes to conduct similar experiments with other kinds of waves, like light or sound, although it may prove more difficult to find an analogue of their time-mirror jolt in other wave mediums, he says.

(1499) "The work is spectacular," says Geoffrey Edelmann of the US Naval Research Laboratory in Washington, D.C. "They essentially change the rules that govern how the wave propagates and, in doing so, create an entirely new way to focus waves free of spatial constraints." He also thinks the technique could be extended to other kinds of waves, and perhaps have medical applications such as lithotripsy, which uses ultrasound waves to break up kidney stones. Journal reference: Nature Physics, DOI: 10.1038/nphys3810

**Time reversal and holography with spacetime transformations Vincent Bacot, Matthieu Labousse, Antonin Eddi, Mathias Fink & Emmanuel Fort Nature Physics (2016) doi: 10.1038/nphys3810**

(1500) Wave control is usually performed by (e) spatially engineering the properties of a medium.

(1501) Because time and space play similar roles (e&eb) in wave propagation, manipulating time boundaries provides (eb) a complementary approach.

(1502) Here, authors experimentally demonstrate the relevance of this concept by (e) introducing instantaneous time mirrors.

(1503) Authors show with water waves that a sudden change of (e) the effective gravity generates (eb) time-reversed waves that refocus at the source.

(1504) They generalize this concept for (e) all kinds of waves, introducing (eb) a universal framework which explains (eb) the effect of any time disruption on (e&eb) wave propagation.

(1505) Authors show that sudden changes of the medium properties generate (eb) instant wave sources that emerge (eb) instantaneously from the entire space at (eb) the time disruption.

(1506) The time-reversed waves originate from (e) these 'Cauchy sources', which are (+) the counterpart of Huygens virtual sources on (eb) a time boundary. It allows us to revisit the holographic method and introduce a new approach for wave control.

**Viewpoint: A Bird's Eye View of Circuit Photons Sebastian Schmidt, Institute for Theoretical Physics, ETH Zurich, 8093 Zurich, Switzerland June 30, 2016• Physics 9, 74 A scanning probe detects the quantum states of photons in a microwave circuit, providing the information needed for quantum simulations. Figure caption expand figure A. Houck/Princeton (deleted)**

(1507) Figure 1: Scanning defect microscopy provides a map of photons in a resonator lattice. Houck and colleagues demonstrated the technique using (e) 49 resonators (grey lines) that were coupled together to (e&eb) form a kagome lattice.

(1508) Quantum mechanics rules the dynamics of (e) light and matter.

(1509) Yet performing a quantum-mechanical simulation of a material from first principles is (=) practically impossible on a classical computer because (e) the complexity of the simulation increases (eb) exponentially with the number of particles involved.

(1510) The solution, according to Richard Feynman, was to build a machine out of (e) quantum building blocks that could directly emulate (e&eb) the material itself [1].

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| (1511) Prototypes of such quantum simulators that are based on (e) ultracold atoms, ions, photons, and superconducting microwave circuits are now available [2], with (e&eb) the latter, in particular, having attracted (e) Silicon Valley's interest. |
| (1512) The challenge with these circuit-based simulators, however, is that they are (=) 2D, which complicates the readout of their constituent elements. Andrew Houck from Princeton University, New Jersey, and colleagues have now delivered an attractive solution by developing a technique [3], called scanning defect microscopy (Fig. 1), that determines (eb) the number of photons occupying (e) each mode of a 2D microwave circuit. |
| (1513) It is this information that would serve as (=) the fundamental input and output for (e) certain quantum simulations. |
| (1514) Superconducting microwave circuits combine electronic and (e&eb) photonic degrees of freedom [4, 5]. |
| (1515) The main element of the circuit is (=) a transmission line, which is made up of (e) a central superconducting wire separated by a gap from (e) two grounded plates. |
| (1516) All of these structures are on a single plane, as if one had taken a 2D slice through (e&eb) a coaxial cable. |
| (1517) When truncated, the transmission line becomes (=) a resonator, which can host (e&eb) discrete photon modes within (eb) its gaps |
| (1518) Large lattices of resonators can be engineered in various 1D or 2D geometries by coupling (e&eb) two, three, or more resonators together via (e&eb) a capacitive interface. In many ways, photons in such devices behave similarly to electrons in a solid. |
| (1519) To make microwave circuits that can simulate (e&eb) quantum phenomena faster than a classical computer, however, resonator lattices have to be integrated with (e&eb) superconducting qubits |
| (1520) . Such qubits are controlled with (e&eb) electrical currents in a Josephson tunnel junction, and in many respects, they behave (e&eb) like artificial atoms, which couple to (e&eb) the photons in the circuit. |
| (1521) As a result, superconducting microwave circuits can be used to (e) explore the coupling between the quantum states of light and (e&eb) matter, the regime of (e) circuit quantum electrodynamics (cQED). |
| (1522) Photons in these devices often exhibit (eb) striking matter-like behavior [6, 7], providing (eb) the basis for the simulation of complex materials. |
| (1523) Such circuits can be fabricated on a substrate using (e) standard lithographic techniques, with (e&eb) qubits and resonators that are (=) hundreds of micrometers or even millimeters in size. |
| (1524) Small-scale realizations of cQED lattices have already been used for (e) proof-of-principle quantum simulations [8, 9]. |
| (1525) However, future large-scale simulations will have to be initialized and measured (eb) efficiently—steps known as (=) drive and readout. |
| (1526) This requires connecting the lattices to (e) external circuitry, such as voltage sources, circulators, and amplifiers, which typically reside outside (e) the chip. |
| (1527) So far, this connectivity has been realized through (e&eb) external transmission lines, but these elements can only couple to (e&eb) the resonators at the edges of (e) the lattice because of (e) its 2D nature. |
| (1528) Such a side view provides (eb) only limited information about (e&eb) the quantum state and spatial correlations of photons inside (eb) the bulk resonators, seriously limiting (e) the applicability of circuit QED devices for (e) quantum simulations. |

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| (1529) The probe developed by Houck and colleagues, scanning defect microscopy (SDM), solves (e) this problem by (e) moving outside the plane of the chip to provide (eb) a fuller image of photons within (eb) the chip. |
| (1530) In SDM, a small sapphire tip is positioned a few tens of (e) micrometers above the surface of one of the resonators (Fig. 1). |
| (1531) The presence of this dielectric modifies (e&eb) the evanescent electric field from the circuit photons, which extends (eb) just outside the plane of the chip. |
| (1532) Specifically, the tip increases (eb) the local capacitance of the transmission line, downshifting (e) the local resonator frequency. |
| (1533) This in turn causes (eb) a drop in the transmission of photons through (e&eb) the lattice when (e) it is driven at its nominal resonance frequency. |
| (1534) Houck's team showed that if the sapphire tip has only a weak (linear) effect on the resonator frequency, then (eb) a simple analytic relationship exists between the induced frequency shift, (e&eb) the change in transmission, and (e&eb) the local photonic occupation of each resonator. |
| (1535) This relationship provides (eb) a recipe for detecting the local photon occupancy inside (eb) the lattice. |
| (1536) First, calibrate the probe by determining (eb) the frequency shift for (e) a single isolated resonator as (=) a function of the tip's position. |
| (1537) Second, measure the transmission through (e&eb) the whole lattice with and without (e&eb) the probe. |
| (1538) Finally, use the measured transmission drop with (e&eb) the probe present to find (eb) the photonic occupation of (e) an individual resonator in (eb) the lattice. |
| (1539) A map of the photon-mode occupancies can be obtained by (e) scanning the probe across the whole lattice. |
| (1540) Houck's team used (e) SDM to image the occupancies of (e) the photon modes of (e) 49 coupled resonators arranged in (eb) a special geometry called a kagome lattice (see details in Fig. 1). |
| (1541) (To simplify the measurement, no superconducting qubits were (=) incorporated with (e&eb) the resonators.) |
| (1542) The occupancies they measured were in excellent agreement with (=) theoretical predictions, corresponding to (e&eb) a fidelity of 0.99. |
| (1543) The high fidelity is even more impressive considering the weak signal-to-noise ratio of (e) their technique (which is limited by (e) the intentionally small size of the frequency-shifted region) and (e&eb) the technical challenge of (e) positioning the probe inside (eb) a cryostat for (e) in situ measurements. |
| (1544) While SDM shares similarities with (e&eb) so-called bead-pull experiments for accelerators and scanning gate microscopy for (e) nanostructures, it is truly a (=) new tool for quantum electronics. |
| (1545) First application of SDM is likely to be (=) the simulation of 2D topological phenomena and (e&eb) geometrically frustrated lattices (of which the kagome lattice is (=) one example.) |
| (1546) Both are associated with (e&eb) macroscopically degenerate systems, which are (=) notoriously difficult to simulate on (e&eb) a classical computer. |
| (1547) Such circuits are (=) genuinely open systems (photons can be added or lost). |
| (1548) This combined with (e&eb) their inherent strong light-matter coupling, makes (eb) such quantum simulators ideal candidates for (e) the study of many-body, out-of-equilibrium phenomena, such as (=) thermalization or nonequilibrium phase transitions. |

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| (1549) SDM could also lead to (eb) the discovery of new exotic phenomena, which are (=) unknown in traditional solid-state systems. |
| (1550) In a nutshell, SDM constitutes a new and indispensible tool for (e) future experiments on quantum simulations with microwave photons. This research is published in Physical Review X. |

**Imaging Photon Lattice States by Scanning Defect Microscopy D. L. Underwood, W. E. Shanks, Andy C. Y. Li, Lamia Ateshian, Jens Koch, and A. A. Houck Phys. Rev. X 6, 021044 – Published 30 June 2016**

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| (1551) Microwave photons inside (eb) lattices of coupled resonators and superconducting qubits can exhibit (eb) surprising matter like behavior. |
| (1552) Realizing such open-system quantum simulators presents (eb) an experimental challenge and requires (e) new tools and measurement techniques. |
| (1553) Here, authors introduce scanning defect microscopy as one such tool and illustrate its use in mapping (e&eb) the normal-mode structure of microwave photons inside (eb) a 49-site kagome lattice of coplanar waveguide resonators |
| (1554) Scanning is accomplished by (e) moving a probe equipped with a sapphire tip across (e&eb) the lattice. |
| (1555) This locally perturbs (e&eb) resonator frequencies and induces (eb) shifts of the lattice resonance frequencies, which authors determine by (e) measuring the transmission spectrum. |
| (1556) From the magnitude of mode shifts, author can reconstruct (eb) photon field amplitudes at (eb) each lattice site and thus create (eb) spatial images of the photon-lattice normal modes. Received 4 December 2015 DOI:http://dx.doi.org/10.1103/PhysRevX.6.021044 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society |

| |
|---|
| |
| We assume that |
| $$D_7 = \mu_7(a_7) + (f_7)_7(P_7)_7 + (f_7)_7(Q)_7 \quad (7)$$ |
| $$n_7 = \beta_7(a_7)h_7(Q_7), \quad (3)$$ |
| Where $\mu_7$the age- dependent debit rate whit is $\mu_7(a_7)$ 0 $\forall a > 0$; $(f_7)_7$ is the density- dependent death rate with $f_7(0) = 0$, $(f_7^7)_7(P_7)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_7)_7(P_7) = (k_7)_7$ ∞; $f_7$ is the functional response of creditors to prey with $[f_7]_7(0) = 0, (f_7^7)_7(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_7)_7(Q_7) = [k_7]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^7(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0. |

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_7 + (\rho_a)_7 = -\mu_7(a_7) + (f_7)_7 P_7) + \big(f_7)_7(Q_7)\big)(\rho_7)_7 \quad (4)$$

$$\rho_7(t_7, 0) = h_7(Q_7) \int_0^\infty \beta_7 (a_7)\rho_7(t, a_7)da, \quad (5)$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(P_7)Q_7 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_7$ and $Q_7$ of (4-6) satisfy

$$0 \leq \rho_7(t, a_7) \leq \rho_7^*(t, a_7),$$

$$0 \leq Q_7(t) \leq Q_7^*(t),$$

Where $\rho_7^*(t, a_7)$ and $Q_7^*(t)$, are solutions of

$$(\rho_t^*)_7 + (\rho_a^*)_7 = -(\mu_7(a_7) + (f_7)_7(P_7^*) + (f_7)_7 Q_7^*))\, \rho_7^*, \quad (7)$$

$$\rho_7^*(t, 0) = \int_0^\infty \beta_7(a_7)\, \rho_7^*(t, a_7)da_7 \quad (8)$$

$$\frac{dQ_7^*}{dt} = -b_7\, Q_7^* + g_7(P_7^*)Q_7^*, \quad (9)$$

With $P_7^*(t) = \int_0^\infty \rho_7^*(t, a_7)da_7$

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From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_7^*(t)}{dt} = P_7^*(t)\big(r_7 - (f_7)(P_7^*) - (f_7)_7(Q^*)\big), \quad (10)$$

$$\frac{dQ_7^*(t)}{dt} = Q_7^*(t)(-b_7 + g_7(P_7^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_7(a)e^{-\int_0^a \mu_7(\tau)dt}e^{-ra}\,da_7 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_7^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(t)dt}\,da_7 > 1, \quad (H1)$$

Which guarantees $r_7 > 0$.

*Remark.* If $\beta_7(a) = [\beta_0]_7 a_7 e^{-\alpha_7 a_7}, \mu_7(a_7) \equiv \mu_7$ then the solution of (12) is

$$r_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_7]_7(P_7) = 0, [f_7](Q_7) = \lambda_7 Q_7\ h_7(Q_7) = 1, \quad and\ \ g_7(P_7) = c_7 P_7$$

$$[\rho_7]_7 + [\rho_a]_7 = -(\mu_7 + \lambda_7 Q_7)\rho_7$$

$$\rho_7(t,0) = \int_0^\infty [\beta_0]_7 a_7 e^{-\alpha_7 a_7} \rho_7(t,a_7)\, da_7$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + c_7 P_7 Q_7$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_7}{dt} = \omega_7 P_7 - \lambda_7 P_7 Q_7$$

$$\frac{dQ_7}{dt} = b_7 Q_7 - c_7 P_7 Q_7$$

Where $\omega_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$

## 24. THE EQUILIBRIUM $(\rho_0)_7(a_7), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_7(a_7), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_7 + (\rho_a)_7 = -\big(\mu_7(a_7) + (f_7)_7(P_7)\big)\rho,$$

$$\rho_7(t,0)$$

$$= \int_0^\infty \beta_7(a_7)\rho_7(t,a_7)da_7 \qquad (13)$$

The solution of (13) has the form of

$$\rho_7(t,a_7) = Ce_7^{r(t-a_7)} - \int_0^a (\tau)d\tau - \int_0^7 (f_7)_7\big(P_7(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_7(t) = C_7 e^{r_7 t - \int_0^t (f_7)_7(P_7(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_7^{-7}(r)$ if $k_7 > r$. This is equivalent to

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{(k_7)_7 a_7} da_7 < 1.$$

In fact, if we define

$$R_7(s_7) = \int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau)d\tau} e^{-s_7 a_7}\, da_7$$

$R(s)$ is a decreasing function. $R_7(r_7) = 1$ And $R_7(k_7)_7 < 1$ imply $r_7 < (k_7)_7$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_7(a_7), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{k_7 a_7} da_7$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_7 \equiv \int_0^\infty (\rho_0)_{77}(a_7)da_7$ *satisfies*

$$\int_0^{\infty} \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{-(f_7)_7(P_0)a_7} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_7(a_7), 0)$

Let $u_7(t, a_7) = \rho(t, a_7) - \rho_0(a_7)$ Then

$$N_7(t) \equiv \int_0^{\infty} u_7(t, a_7) da_7 = \int_0^{\infty} \rho_7(t, a_7) da_7 - \int_0^{\infty} (\rho_0)_7(a_7) da = P_7(t) - P_0)_7$$

And

$$(u_7)_7 + (u_a)_7 = (\rho_t)_7 + (\rho_a)_7 - (\rho_0^7)_7(a_7)$$

$$= -\mu_7(a_7) + (f_7)_7(P_7) + (f_7)_7(Q_7))(u_7 + (\rho_0)_7) + \mu_7(a_7)(\rho_0)_7 + (f_7)_7(P_0)_7(\rho_0)_7$$

$$= -\mu_7(a_7)u_7 - (f_7)_7(P_7) + (f_7)_7(Q_7))\, u_7 - (\rho_0)_7\, (f_7)_7(P_7) - (f_7)_7(P_0)_7$$

$$+ (f_7)_7(Q_7)), \qquad (15)$$

Where $(\rho_0^7)_7(a_7)$ denotes $d(\rho_0)_7(a_7)/da_7$

Expansion $(f_7)_7(P_7)$ and $(f_7)_7(Q_7)$ at $(P_0)_7$ and 0 yield respectively

$$(f_7)_7(P_7) = (f_7)_7((P_0)_7 + N_7)$$

$$= (f_7)_7(P_0)_7 + (f_7^7)_7(P_0)_7 N_7 + o_7(N_7), \qquad (16)$$

$$(f_7)_7(Q_7)$$

$$= (f_7^7)_7(0)Q + o_7(Q_7) \qquad (17)$$

Write $f_7(P_0) = f_{70}, f_7'(P_0) = f_{70}'$, and $f_7'(0) = f_{70}'$ and substitute (16) and (17) into (15) to obtain

$$(u_7)_7 + (u_a)_7 \approx -(\mu_7(a_7) + (f_{70})_7)u - (f_{70}^7)_7(\rho_0)_7(N_7) - (f_{70}^7)_7(\rho_0)_7(Q_7)$$

$$u_7(t, 0) = \rho_7(t, 0)(\rho_0)_7(0)$$

$$= h_7(Q_7)\int_0^{\infty} \beta_7(a_7)(u_7(t,a) + (\rho_0)_7(a_7)da_7 - (\rho_0)_7\,(0)$$

$$= h_7(Q_7)\int_0^{\infty} \beta_7(a_7)(u_7(t,a_7)da_7 + (h_7(Q_7) - 1)(\rho_0)_7\,(0) \qquad (18)$$

Expanding

$$h_7(Q_7) = 1 + \ (h^7)_7(0)Q_7 + 0(Q_7) \equiv 1 + h_0^{'}Q_7 + 0Q_7$$

And substituting it into (18), we arrive at

$$u_7(t,0) \approx \int_0^{\infty} \beta_7\,(a_7)u_7(t,a_7)da_7 + (h_{a_7}^{'})(\rho_0)_7Q_7$$





The linearization is the following :

$$[u_7] + [u_7] = -(\mu_7(a_7) + [t_{70}]_7u_7 - [\alpha_7]_7(a_7)N_7 - [\alpha_7]_7(a_7)Q_7$$

$$u_7(t,0) = \int_0^{\infty} \beta_7(a_7)u_7(t,a_7)da_7 + [\xi_7]Q_7 \qquad (19)$$

$$\frac{dQ_7}{dt} = -b_7Q_7$$

where

$$[\alpha_7]_7(a_7) = [t_7^7a]_7[p,o]_7(a_7)d2(a_7) = [f_{70}^7]_7(\rho\ o\ )_7(a_7)\xi = [\mathrm{h\ o}^7]_7[\rho\mathrm{o}]_7(\mathrm{o}) \text{ and } \mathrm{b}_7 = \mathrm{b}_7 - g_7(P_7)$$

If $g_7(P_0) > b$, clearly, $(\rho_0)_7(a_7), 0)$ is unstable and the predator population is persistent. Assume $g_7(P_0)_7 < b$.

Then, $Q_7(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_7 + u_a = -(\mu_7(a_7) + (f_{70})_7)u_7 - (\alpha_7)_7(a_7)N_7$$

$$= -(\mu_7(a_7) + f_{70})u_7 - \alpha_7(a_7)\int_0^\infty u_7(t, a_7)da_7$$

$$u_7(t, 0) = \int_0^\infty \beta_7\,(a_7)u_7(t, a_7)da_7$$

Using the method developed by Cushing [2, 3], we suppose

$$u_7(t, a_7) = \tilde{u}_7(a_7)^{e^{\gamma(t-a_7)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_7}\,da_7$. Then,

A solution of the first differential equation for $\tilde{u}(a_7)$ gives,

$$\tilde{u}_7(a_7) = e^{-\int_0^a \mu_7(\tau_7)d\tau - (f_{70})_7 a}\left(\tilde{u}_7(0) - w\int_0^a e^{\int_0^\tau \mu_7(s_7)ds + (f_{70})}\tau_7\,(\alpha_7)_7(\tau_7)e^{\gamma_7\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_7(a_7)e^{-\gamma_7 a_7 - (f_{70})_7 a_7 - \int_0^a \mu_7(\tau_7)d\tau}\right)\tilde{u}_7(0)$$

$$+w\int_0^\infty \beta_7(a_7)e^{-\gamma_7 a_7 - (f_{70})_7 a_7 - \int_0^a \mu_7(\tau_7)d\tau}\int_0^a (\alpha_7)_7(\tau_7)e^{-\gamma_7\tau_7}e^{(f_{70})_7\tau_7 + \int_0^\tau \mu_7(s_7)ds}d\tau_7\,da_7$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty} e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau} da\right)(\tilde{u}_7(0))_7$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau}\int_0^a e^{(f_{70})_7\tau_7+\int_0^\tau \mu_7(s_7)ds+y\tau}\alpha_7(\tau)d\tau\, da\right)w=0. \quad (23)$$

Define $Y(\gamma,a)=e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(t_7)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)da_7\right)u_7(0)$$

$$+w\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\, da_7=0, \quad (24)$$

$$\left(\int_0^{\infty}Y_7(\gamma_7,a_7)da_7\right)\tilde{u}_7(0)$$

$$+\left(-1-\int_0^{\infty}Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\, da_7\right)w=0.$$

Equation (20) has a solution of the form of $u_7(t,a_7)=u_7(a_7)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_7(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_7(\gamma_7,a_7)=\int_0^a(\rho_0)_7(\tau_7)Y^{-7}(\gamma_7,a_7)d\tau_7$. We obtain the characteristic equation

$$f_{70}^{l}{}_{7}\left(\int_0^{\infty} Y_7(\gamma_7,a_7)da_7\right)\left(\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$+\left(1-\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)da\right)\left(1+(f_{70}^{l})_7\int_0^{\infty}Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_7(a_7)=(\rho_0)_7(0)e^{-\int_0^a \mu_7(\tau_7)d\tau_7-(f_{70})_7a_7}$$

$$I_7(\gamma_7,a_7)=\int_0^a(\rho_0)_7(\tau_7)e^{\gamma_7\tau_7+(f_{70})_7\tau+\int_0^{\tau}\mu_7(s_7)ds_7}d\tau$$

$$=\frac{(\rho_0)_7(0)}{\gamma_7}(e^{\gamma_7a_7}-1)$$

Let

$$E_7(a_7)=e^{-(f_{70})_7a_7-\int_0^a\mu_7(\tau_7)d\tau_7}$$

Then $[Y_7]_7[\gamma_7]_7,a_7)=E_7(a_7)e^{-\gamma_7a_7}$. by a simple calculation,

$$\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da=\frac{[\rho_0]_7(0)}{\gamma}\left(1-\int_0^{\infty}\beta_7(a_7)E_7(a_7)e^{-\gamma_7a_7}\ da\right);$$

$$\int_0^{\infty}Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da=\frac{[\rho_0]_7(0)}{\gamma_7}\left(\int_0^{\infty}E_7(a_7)e^{-\gamma_7a_7}\ da_7\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_7(0)}{\gamma_7}[f_{70}^7]\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\,(\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7)$$

$$+\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right).$$

$$\left(1+\frac{(\rho_0)_7(0)}{\gamma_7}(f_{70}^{'})_7\left(\int_0^\infty E_7(a_7)da_7-\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right)(1+\frac{(f_{70}^{'})_7}{\gamma_7}(\rho_0)_7(0)\int_0^\infty E_7(a_7)da_7)=0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)da_7 = 1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7 \neq 0,$$

Since $\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)\int_0^a \alpha_7(\tau_7)Y^{-7}(\gamma_7,\tau_7)d\tau_7 da_7 > 0$. Therefore, the characteristic solution is

$$\gamma_7 = -(f_{70}^{'})_7(\rho_0)_7\int_0^\infty E_7(a_7)da_7.$$

Since $f_7'(P_0)_7 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_7), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE OF DEBIT AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_7), 0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_7 + (\mu_7(a_7) + [f_7]_7(P_7^*) + [f_7]_7(Q_7^*))P_7^* = 0, \quad (76)$$

$$p_7^*(0) = h_7(Q_7^*)\int_0^\infty \beta_7(a_7)p_7^*(a_J)p_7^*(a_7)da_7 \quad (77)$$

$$g_7(P_7^*) - b_7 = 0, \quad (78)$$

Where $P_7^* = \int_0^\infty p_7^*(a_7)\, da_7$

Equation (28) yields $P^* = g^{-7}(b)$,

From (26),

$$p_7^*(a) = p_7^*(0)e^{\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(Q_7^*)a_7}$$

Substituting into (27), we need to solve

$$\left(1-h_7(Q_7^*)\int_0^{\infty}\beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7-[f_7]_7(P_7^*)a_7-[f_7]_7(Q_7^*)a_7}da_7\right)p_7^*(0)=0$$

For $Q_7^*$

Fix $p_7^*(0)$ and $P_7^*$ Let

$$H_7(z_7)=1-h_7(z_7)\int_0^{\infty}\beta_7(a_7)\,e^{-\int_0^a \mu_7(\tau_7)d\tau_7-[f_7]_7(P_7^*)a_7-[f_7]_7(z_7)a_7}da_{,7}$$

And assume that

$$[k_3]_7\int_0^{\infty}\beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7-[f_7]_7(g_7^{-7}(b_7))a_7-[k_7]_7a_7}\,da_7<1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_7(z_7)]>0.$$

Clearly,

$$\frac{dH_7}{dz_7}=h_7(z)(\,[f_7^{'}]_7(z_7)\,)\int_0^{\infty}a_7\beta_7(a_7)\,e^{-\int_0^a \mu_7(\tau_7)d\tau_7-f_7(P_7^*)a_7-f_7(z_7)a_7}da_{7}$$

$-h^{'}(z_7)\int_0^{\infty}\beta_7(a_7)\,e^{\int_0^a \mu_7(\tau_7)d\tau_7-[f_7]_7(P_7^*)a_7-[f_7]_7(z_7)a_7}\quad da_7$

(See author's It is a debit-credit world)

By the assumptions on $\left[f_7^7\right]_7 and\ [h^7]^7 dH_7/dz > 0$. Hence, if

$$H_7(0) = 1 - \int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7} da_7 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_7(0) < 0$ follows if $g_7^-(b_7) < [f_7^{-7}]_7(r_7)$, where $r_7$ is the solution of (12).

Indeed, since

$$R_7(s_7) = \int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 e^{-y\frac{a}{7}}} da_7$$

Is a decreasing function of $s_7$ when $P_7^* = [g^{-7}]_7(b_7) < [f_7^{-7}]_7(r_7)$, that is $[f_7]_7(P_7^*) < r$, $R\left(f_7(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

**THEOREM 3**

Assume that the hypothesis (H3) holds and $R(f_7(g^{-7}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7\left(g_7^{-7}(b)\right)a_7} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_7^*(a_7), Q_7^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_7^*(a_7), Q_7^*)$ is more complicated. We again consider perturbations $m_7(t_7 a_7) = p_7(t_7, a_7) - p_7^*(a_7)$ and $W_7(t) = Q_7(t) - Q^*$.Then

$M_7(t) \equiv \int_0^\infty m_7(t,a_7)da_7 = P_7(t) - P_7^*,$

$$[m_7]_7 + [m_a]_7 = [p_7]_7 + [p_a]_7 - [p^*]_7(a_7) \approx -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7$$

$$-[f_{7*}^7]_7 p_7^*(a_7)M - [f_{7*}^7]_7 p_7^*(\mathrm{a})W_7,$$

$$m_7(t,0)p_7(t,0) - p_7^*(0) \approx h_* \int_0^\infty \beta_7(a_7)m_7(t,a_7)\,da + \frac{[h_*^7]_7}{[h_*]_7}\, p_7^*(0)W_7$$

$$\frac{dW_7}{dt} = \frac{dQ_7}{dt} \approx [g_*^7]_7 Q_7^* M,_7$$

Where $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}']_7$ denoting $[f'_7(P^*)]_7[f_{7*}']_7$ denoting

$[f_7]_7$ $(Q_7^*,$ $h_{7}^*$ denoting $h_*'$ denoting $h_7'$ $(Q_7^*)$, and $[g_*']_7$

Denoting $[g']_7(P_7^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$[m_7]_7 + [m_a]_7 = -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7 - [f'_{7*}]_7[p_7^*](a_7)M_7 - [f_{7*}']_7 p_7^*(a_7)W_7$

$m_7(t,0) = [h_*]_7 \int_0^\infty \beta_7(a_7)m_7(t,a_7)da_7 + \frac{[h_*^7]_7}{[h^*]_7}[p^*]_7(0)\,W,_7$ (30)

$$\frac{dW_7}{dt} = [g_*^7]_7[Q_7^*]M_7$$

Integrating $[m_7]_7(t,a_7)$ with respect to a from 0 to ∞, we have

$$\frac{dM_7(t)}{dt} = \int_0^{\infty} [mt]_7(t, a_7)da_7$$

$$= -\int_0^{\infty} [m_a]_7\,(t, a_7)da_7 - \int_0^{\infty}\ \mu(a_7)m(t, a_7)da_7$$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f^{'}{}_{7*}]_7 \int_0^{\infty} p_7^*(a_7)da_7\right)M_7 - [f_{7*}^7]_7 \int_0^{\infty} p_7^*(a_7)da_7 W_7$$

$$= m_7(t, 0) - \int_0^{\infty} \mu_7(a_7)\, m_7(t, a_7)da_7$$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 - [f_{7*}^{'}]_7 P_7^* W_{.7}$$

From (30), we arrive at the equations

$$\frac{dW_7(t)}{dt} = \int_0^{\infty} ([h_*]_7\beta_7(a_7) - \mu_7(a_7))m_7(t, a_7)da_7$$

$$-\left([f_{7*}] + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 + \left(\frac{[h_*^{'}]_7}{[h_*]^7}\ P_7^*(0) - f_{7*}^{'} P_7^*\right)W_7,$$

(31)

$$\frac{dW_7(t)}{dt}[g^{'}{}_*]_7\, Q_7^*\, M_{.7}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{"*}]_7\beta_7(a_7) - \mu_7(a_7))m_7(t, a_7)da_7 = ([h_*]_7\beta_7(\theta_7) - \mu_7(\theta_7))M,_7$$

Where $0< \theta_7 = \theta_7(t) < \infty$.

And if a $_7^7$ [a $_7^7$]2<4 [a $_7$]2<4 a $_7$e,

Re $\lambda_7[\lambda_7]_7 = \mathrm{Re}\lambda_7[\lambda_7]_7 = -\frac{7}{7}a_7[\mathrm{a}2\ ]_7 < 0$.

Let ⊿=max$_{7<u<\infty}$ |h $_*$ [h $_*$] $_7$β $_7$($a_7$) − $\mu_7(a_7)$|. Then|$n_7$(t)| ⊿∀t . Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium(ρ $^*$ (a), $\boldsymbol{Q}^*$) is asymptotically stable if either a $_7^7$ <4 a $_7$ e and [a $_7$] > 1⊿,or [a $_7^7$] [a $_7$]e and (a $_7$ − $\sqrt{(\mathrm{a}_7^7)_7 - 4(\mathrm{a}_7)_7\mathrm{e})} > 1⊿$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{\mathrm{a}}{\mathrm{a}+7}\right)$ and β $_7$($a_7$) = ae-$^{\mathrm{ma}}$ Then ⊿=[$\mu_0$] + 2 Simple calculation leads to

$$\int_0^\infty \beta_7(a_7)\mathrm{e}^{-\int_0^{\mathrm{a}}\mu_7(\tau_7)+d\tau_7\mathrm{da}} = \frac{1}{(\mathrm{m}+\mu_0+1)^7} + \frac{1}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_7$ + [$\mu_0$] .5. Assume[$f_7$]($x_7$)= $K_7x_7$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_7(x_7)$= [0$x_7$]/(2 + $x_7$); the functional responses [[$f_7$]($x$) = [0] $x_7$/(1 + $x_7$)and $h_7(x_7)$=(1+$\omega_7x_7$)/(1 + $x_7$)(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_7 + [\mu_0] + [\mathrm{K}][P_0])}{[m_7]_7 + (\mu_0)_7 + [\mathrm{K}][P_0])_7^3\ [\mathrm{K}]_7[P_0]_7^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1+P_0) < b_7$ i e $[P_0] < b/(\delta_7 - b_7)$then the equilibrium$(P_0)$ $[P_0]_7(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_7$

$[f_{7*}] = K_7\ \mathrm{P}^*\ [f_{7*}] = \frac{\sigma \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*} \quad [h_*] = \frac{7+\omega_7 \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*}$

$[f_{7*}] = K_7\ [\mathrm{f}_{7*}^{\ 7}] = \frac{\sigma_7}{(7+\mathrm{Q}_7^*)^7} [\ \mathrm{g}_*^{\ 7}] = \frac{\sigma_7}{(7+\mathrm{P}_7^*)^7} \quad [\mathrm{h}_*^{\ 7}] = \frac{\omega_7 - 7}{(7+\mathrm{Q}_7^*)^7}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_7}{(m + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma)^7} + \frac{1\omega_7}{([m_7]_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma_7)^3} < 1$$

And

$$\frac{1}{(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^7} + \frac{1\omega_7}{(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $\mathrm{Q}^*$ which can be solved from the following equation:

$(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7} + (m_7 + [\mu_0]_7 + \frac{K_7 b_7}{\delta_7 - b_7} + \sigma_7)\ \mathrm{Q}_7^*)^3$

$= m_7 + [\mu_0]_7 + 3 + \frac{K_7 b_7}{\delta_7 - b_7}(1 + \mathrm{Q}^*)(1 + \omega_7 \mathrm{Q}_7^*$

Now

$$\left[\mathrm{a}\begin{matrix}1\\1\end{matrix}\right]-4\left[\mathrm{a}_{7}\right]\mathrm{e}=\left(\frac{1K_7 b_7}{\delta_7-b_7}\right)^{7}+\frac{4K_7 b_7\sigma_7\,\mathrm{Q}_{7}^{*}}{(\delta_7-b_7)(7+\mathrm{Q}_{7}^{*})}+\frac{\sigma_{7}^{7}\mathrm{Q}_{7}^{*7}}{\left(1+\mathrm{Q}_{7}^{*}\right)^{7}}$$

$$\frac{4\,b_7\sigma_7\,(\delta_{7}-\mathrm{b})\mathrm{Q}^{*}}{\delta_{7}\,(1+\mathrm{Q}_{7}^{*})^{7}}-\frac{4\,(\delta_{7}-b_7)^{7}\,(1-\omega_7)\mathrm{Q}_{7}^{*}}{\delta_{7}\left(1+\omega_7\mathrm{Q}_{7}^{*}\right)(1+\mathrm{Q}_{7}^{*})}\mathrm{P}*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_{7}\,ae_{7}^{mu}\;\;[\mu_0]_{7}-\frac{a_7}{a_7+7}$$

Hence if $[\mathrm{a}_{7}^{7}>4(a_7)_{7}\mathrm{e}$ and

In the following are given some more modular structures that are applied with the model and the concomitant results derived thereof:

Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments

(1) We have an intuitive understanding of how the world works. Our intuition works well in our everyday lives, but it doesn't quite match (e&eb) reality.

(2) One of the basic assumptions authors make about objects is that they are (=) objectively real.

(3) For example, when you leave your apartment in the morning, you imagine your house is right where you left it as you go about your day. Barring some calamity, you expect to find your apartment just as you left it when (e) you return in the evening.

(4) This idea is known as Macrorealism, in that a macroscopic object such as your house is still there even when (e) you're not looking at it

(5) In scientific terms we would say the state of an object exists independent of (e) an observer. While this seems obviously true, we can test the idea experimentally.

(6) To show this, let's consider an experiment. Suppose we have a box that contains a coin. At any time we can open the box and (eb) see whether it's heads or tails.

(7) The catch is (=) that the box has to follow a certain rule. Whatever you observe initially, (eb) you must observe the same thing ten seconds later.

(8) If you open the box and see heads, then (eb) ten seconds later when you open the box you must also see heads.

$$[a_7] = \frac{1K_7 b_7}{\delta_7 - b_7} + \frac{\sigma_7 \, \mathrm{Q}_7^*}{1 + \mathrm{Q}_7^*} > 1\vartriangle$$

This equation is important in its consequence in respect of following cases for which the model is applied:

**Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments**

(13) For the experiment we'll open the box, then open it five seconds later, then finally open it again at ten seconds. We know the first and last observations have to agree, but what about the observation in the middle? It depends upon what the box does with (e&eb) the coin.

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| (14) One possibility is that the box does nothing, so if you see tails the first time (eb), you will see tails both other times. |
| (15) Another possibility is that each time you open the box (eb) it flips the coin when you close it. |
| (16) So the first time you see heads, (eb) the second time tails, and (eb) the third time heads again. |
| (17) But they could also make a tricky box that flips (e&eb) the coin at seemingly random times, but does so in such a way that every ten seconds it repeats (eb) the pattern. In that case, authors would always see the same thing the first and last time, but the middle observation would seem (eb) random. |
| (18) If Macrorealism is true, then (eb) middle observations being always the same as (=) the first and last observation, always opposite or always random are (=) the only possibilities for the experiment |

**Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry Download as PDF (Size: 62KB) HTML PP. 50-53 DOI: 10.4236/jqis.2011.12007 5,508 Downloads 14,712 Views Citations Author(s) Leave a comment Mohamed S. El Naschie**

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| (13) Building upon the pioneering work of J. Bell [1] and an incredible result due to L. Hardy [2] it was shown (eb) that the probability of quantum entanglement of two particles is (=) a maximum of 9.0169945 percent [2]. |
| (14) This happens to be exactly the golden mean to (e) the power of five (?5) [3-7]. |
| (15) Although it has gone largely unnoticed for a long time, this result was essentially established independently in a much wider context by (e) the present author almost two decades ago [3-6]. |
| (16) The present work gives two fundamentally different derivations of (e) Hardy's beautiful result leading to (eb) precisely the same general conclusion, namely that by virtue of (e) the zero measure of the underlying |
| (17) Cantorian-fractal spacetime geometry the notion of (e) spatial separability in quantum physics is (=) devoid of any meaning [7]. The first derivation is purely logical and uses (e) a probability theory which combines the discrete with (e&eb) the continuum. |
| (18) The second derivation is purely geometrical and topological using (e) the fundamental equations of a theory developed by the author and his collaborators frequently referred to as E-infinity or Cantorian spacetime theory [3-7]. KEYWORDS Hardy's Quantum Entanglement, Golden mean, Cantor sets, Fractal Spacetime, E-Infinity Theory, Quantum Mechanics, J. S. Bell Cite this paper M. El Naschie, "Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry," Journal of Quantum Information Science, Vol. 1 No. 2, 2011, pp. 50-53. Doi: 10.4236/jqis.2011.12007. |

**Chaos, Solitons & Fractals Volume 28, Issue 4, May 2006, Pages 857–878 Sir Hermann Bondi 1919–2005 Cantorian spacetime and Hilbert space: Part I—Foundations G. Iovane http://dx.doi.org/10.1016/j.chaos.2005.08.074**

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| (21) Authors are going to show the link between the ε(∞) Cantorian space and( e&eb) the Hilbert spaces H(∞). |
| (22) In particular, El Naschie's ε (∞) is (=) a physical spacetime, i.e. an infinite dimensional fractal space, where (e) time is spacialized and the transfinite nature manifests (eb) itself. |
| (23) El Naschie's Cantorian spacetime is an arena where (e) the physics laws appear at each scale in a self-similar way linked to (e&eb) the resolution of the act of observation. |
| (24) By contrast the Hilbert space H (∞) is (=) a mathematical support, which describes (eb) the interaction between the observer and (e&eb) the dynamical system under (e&eb) measurement. |

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| (25) The present formulation, which is based on (e) the non-classical Cantorian geometry and topology of spacetime, automatically solves (e) the paradoxical outcome of (e) the two-slit experiment and the so-called particle–wave duality. |
| (26) In particular, measurement (i.e. the observation) is equivalent to (=) a projection of ε(∞) in the Hilbert space built on 3 + 1 Euclidean spacetime. |
| (27) Consequently, the wave–particle duality becomes (=) a mere natural consequence of (e) conducting an experiment in (eb) a spacetime with (e&eb) non-classical topological and (e&eb) geometrical structures, while (e) observing and taking measurements in (eb) a classical smooth 3 + 1 Euclidean spacetime. |
| (28) In other words, the experimental fact that a wave–particle duality exists is (=) an indirect confirmation of the existence of ε(∞) and a property of (e) the quantum–classical interface. |
| (29) Another direct consequence of the fact that real spacetime is) =) infinite dimensional hierarchical ε(∞) is (=) the existence of scaling law R(N), introduced by the author, which generalizes the Compton wavelength. |
| (30) It gives an answer to the problem of (e) segregation of matter at different scales, and shows (eb) the role of fundamental constants such as (=) the speed of light and Plank's constant h in the fundamental lengths scale without (e) invoking the principles of quantum mechanics. Copyright © 2005 Elsevier Ltd. All rights reserved |

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**Chaos, Solitons & Fractals Volume 14, Issue 9, December 2002, Pages 1341–1351 Fractal strings as an alternative justification for El Naschie's Cantorian spacetime and the fine structure constant Carlos Castro, http://dx.doi.org/10.1016/S0960-0779(02)00110-8**

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| (17) Beginning with the most general fractal strings/sprays construction recently expounded in the book by Lapidus and Frankenhuysen, it is shown how the complexified extension of El Naschie's cantorian-fractal spacetime model belongs to (e&eb) a very special class of families of fractal strings/sprays whose scaling ratios are given by (e) suitable pinary (pinary, p prime) powers of the Golden Mean. |
| (18) They then proceed to show why the logarithmic periodicity laws in (eb) Nature are direct physical consequences of (e) the complex dimensions associated with (e&eb) these fractal strings/sprays. |
| (19) Authors proceed with a discussion on quasi-crystals with (e&eb) p-adic internal symmetries, von Neumann's Continuous Geometry, the role of wild topology in (e&eb) fractal strings/sprays, (e&eb) the Banach-Tarski paradox, and tesselations of (e) the hyperbolic plane, quark confinement and (e&eb) the Mersenne-prime hierarchy of (e) bit-string physics in determining the fundamental physical constants in Nature. Fax: +1-914-965-3397 Copyright © 2002 Published by Elsevier Ltd. About ScienceDirectRemote access Shopping cart Contact and support Terms and conditions Privacy policy Cookies are used by this site. For more information, visit the cookies page. Copyright © 2016 Elsevier B.V. or its licensors or contributors. Science Direct ® is a registered trademark of Elsevier B.V. |

**Chaos, Solitons & Fractals Volume 21, Issue 2, July 2004, Pages 261–271 On Cantorian spacetime over number systems with division by zero Jakub Czajko, http://dx.doi.org/10.1016/j.chaos.2003.12.046**

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| (20) Division by zero is proposed in order to open (eb) new horizons for (e) abstract mathematics and physics. |
| (21) Although the formerly forbidden operation may be of (e) little practical value, the feasibility of unrestricted operations is needed for (e) meaningful deployment of linear vector spaces, whose duals are (=) inverse transposes of their primary vector spaces. |
| (22) Both quantum mechanics and (e&eb) pan-geometry of multispatial hyperspace operate on (e&eb) vectors that belong to (e&eb) mutually dual linear vector spaces. |
| (23) The division by zero also may remove (e) many years old paradox related to (e&eb) multiplication of zeros and makes it possible to introduce (eb) the concept of bigroup as (=) an additive and multiplicative group at the same time. |
| (24) This operational extension to (e&eb) algebraic groups makes thus (eb) an infinite-dimensional, fractal Cantorian spacetime (=) the preferred, invertible, operationally unrestricted, abstract mathematical infrastructure for physics. Tel.: +1-925-681-2058 Copyright © 2003 Elsevier Science Ltd. All rights reserved. |

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The nontrivial equilibrium is stable.

Some of the examples in which stability analysis is of seminal importance are: These are the system to which the

model becomes applicable

**Observation of a Dissipation-Induced Classical to Quantum Transition J. Raftery, D. Sadri, S. Schmidt, H. E. Türeci, and A. A. Houck Phys. Rev. X 4, 031043 – Published 8 September 2014**

(279)Here, authors report the experimental observation of (e) a dynamical quantum phase transition in (eb) a strongly interacting open photonic system.

(280) The system studied, comprising a Jaynes-Cummings dimer realized on (eb) a superconducting circuit platform, exhibits (eb) a dissipation-driven localization transition.

(281)Signatures of the transition in the homodyne signal and photon number reveal (eb) this transition to be from (e) a regime of classical oscillations into (e&eb) a macroscopically self-trapped state manifesting revivals, (=) a fundamentally quantum phenomenon.

(282) This experiment also demonstrates (eb) a small-scale realization of a new class of quantum simulator, whose well-controlled (e&eb) coherent and dissipative dynamics is suited to (e) the study of quantum many-body phenomena out of (e) equilibrium. Received 10 December 2013 DOI:http://dx.doi.org/10.1103/PhysRevX.4.031043 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**Viewpoint: Fermionic Vortices Find their Dual John McGreevy, Department of Physics, University of California at San Diego, La Jolla, CA 92093-0319, USA June 27, 2016• Physics 9, 72Theoretical work reveals a surprising relationship between the physics of fermionic vortices and quantum electrodynamics. Figure caption expand figure APS/Alan Stonebraker Figure 1: The action of discrete symmetries on the fermionic particle (top) and its dual fermionic vortex (bottom).**

(283)Dirac fermions in 2+1 dimensions (two spatial dimensions plus time) have been found in (eb) graphene and on (e&eb) the surfaces of topological insulators (TIs).

(284) A trio of papers now shows that these particles can also be found in (eb) a dual description of quantum electrodynamics (QED) [1–3].

(285)This deceptively simple statement joins (e&eb) together several disparate lines of progress in theoretical physics: interacting topological insulators, (e&eb) anomalies in quantum field theory, (e&eb) spin liquids, and (e&eb) compressible Quantum Hall States.

(286)Humans are not very good at understanding interacting many-body systems.

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| (287)We understand the behavior of free particles and can sometimes use (e) perturbation theory to (e) incorporate weak interactions between (e&eb) them. |
| (288)But what happens when we couple (e&eb) the particles strongly? |
| (289)To reveal the interesting collective phenomena that can result, we've had to rely largely on (e) experiments. |
| (290)Examples include (e) superconductivity (e&eb) superfluidity (e&eb) the integer and (e&eb) fractional quantum Hall effects (e&eb) the fractionalization of charge and spin, and (e&eb) the wetness of water. |
| (291)The fractional quantum Hall effect (FQHE) is a physical phenomenon in which the Hall conductance of 2D electrons shows (eb) precisely quantised plateaus at (eb) fractional values of (e) {\displaystyle e^{2}/h} e^{2}/h. |
| (292) It is a property of (e) a collective state in which electrons bind (e&eb) magnetic flux lines to make (eb) new quasiparticles, and excitations have (e) a fractional elementary charge and possibly also fractional statistics. The 1998 Nobel Prize in Physics was awarded to Robert Laughlin, Horst Störmer, and Daniel Tsui for the discovery and explanation of the fractional Hall effect.[1] |
| (293) However, Laughlin's explanation was (=) a phenomenological guess and only applies to (e&eb) fillings {\displaystyle \nu =1/m} \nu =1/m where (e) {\displaystyle m} m is an odd integer. The microscopic origin of the FQHE remains unknown and is a major research topic in condensed matter physics. (Wikipedia) |
| (294)The quantum Hall effect (or integer quantum Hall effect) is a quantum-mechanical version of (e) the Hall effect, observed in (eb) two-dimensional electron systems subjected to (e&eb) low temperatures and strong magnetic fields, in which the Hall conductance σ undergoes (e&eb) certain quantum Hall transitions to take on (e&eb) the quantized values {\displaystyle \sigma ={\frac {I_{\text{channel}}}{V_{\text{Hall}}}}=\nu \;{\frac {e^{2}}{h}},} \sigma = \frac{I_\text{channel}}{V_\text{Hall}} = \nu \; \frac{e^2}{h}, where {\displaystyle I_{\text{channel}}} I_\text{channel} is the channel current, {\displaystyle V_{\text{Hall}}} V_\text{Hall} is the Hall voltage, e is the elementary charge and h is Planck's constant. The prefactor ν is known as the "filling factor", and can take on either integer (ν = 1, 2, 3, ...) or fractional (ν = 1/3, 2/5, 3/7, 2/3, 3/5, 1/5, 2/9, 3/13, 5/2, 12/5, ...) values. |
| (295) The quantum Hall Effect is referred to as (=) the integer or fractional quantum Hall effect depending on whether ν is an integer or fraction, respectively. |
| (296)The striking feature of the integer quantum Hall Effect is (=) the persistence of the quantization (i.e. the Hall plateau) as (e) the electron density is varied. |
| (297) Since the latter remains constant when (e) the Fermi level is in a clean spectral gap, it must be that the Fermi level is immersed in (eb) essential spectrum (i.e. of finite density of states), though Anderson localized. |
| (298) The quantization and stability of the Hall conductance in such conditions was achieved in 1994 by (e) Bellissard, van Elst and Schulz-Baldes using the methods of (e) Non-Commutative Geometry.[1] |
| (299)The fractional quantum Hall effect is more complicated, as its existence relies fundamentally on (e) electron–electron interactions. |
| (300)Although the microscopic origins of the fractional quantum Hall effect are unknown, there are several phenomenological approaches that provide (eb) accurate approximations. |
| (301) For example the effect can be thought of as (=) an integer quantum Hall effect, not of electrons but of (e) charge-flux composites known as (=) composite fermions. |

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| (302) In 1988, it was proposed that there was quantum Hall effect without Landau levels. This quantum Hall Effect is referred to as (=) the quantum anomalous Hall (QAH) effect. |
| (303)There is also a new concept of the quantum spin Hall effect which is an analogue of the quantum Hall effect, where (e) spin currents flow instead of (e&eb) charge currents.[2] |
| (304)In two dimensions, when classical electrons are subjected to (e&eb) a magnetic field they follow (eb) circular cyclotron orbits. |
| (305)When the system is treated quantum mechanically (eb) these orbits are quantized. |
| (306)The energy levels of these quantized orbitals take on (e) discrete values |
| (307):{\displaystyle \textstyle E_{n}=\hbar \omega _{c}(n+1/2),} \textstyle E_{n}=\hbar \omega _{c}(n+1/2), where ωc = eB/m is the cyclotron frequency |
| (308) These orbitals are known as Landau levels, and at weak magnetic fields, their existence gives rise to (eb) many interesting "quantum oscillations" such as the Shubnikov–de Haas oscillations and the de Haas–van Alphen effect (which is often used to map the Fermi surface of metals). |
| (309) For strong magnetic fields, each Landau level is (=) highly degenerate (i.e. there are many single particle states which have the same energy En). |
| (310)Specifically, for a sample of area A, in magnetic field B, the degeneracy of each Landau level is (=) {\displaystyle \textstyle N=g_{s} BA/\phi _ {0},} \textstyle N=g_{s} BA/\phi _ {0}, where gs represents (eb) a factor of 2 for spin degeneracy, and ϕ0 = 2·10−15 Wb is the magnetic flux quantum. |
| (311) For sufficiently strong B-fields, each Landau level may have (e) so many states that all of the free electrons in the system sit in only a few Landau levels; it is in this regime where (e) one observes the quantum Hall Effect. |
| (312)The integers that appear in the Hall Effect are (=) examples of **topological quantum numbers**. |
| (313) They are known in mathematics as the first **Chern numbers** and are closely related to (e&eb) **Berry's phase.** |
| (314)A striking model of much interest in this context is the Azbel-Harper-Hofstadter model whose quantum phase diagram is (=) the Hofstadter butterfly shown in the figure. (deleted due to spatial constraints) |
| (315)The vertical axis is (=) the strength of the magnetic field and the horizontal axis is the chemical potential, which fixes the electron density. The colors represent the integer Hall conductance. Warm colors represent positive integers and cold colors negative integers. The phase diagram is fractal and has structure on all scales. In the figure there is an obvious self-similarity. |
| (316)Concerning physical mechanisms, impurities and/or particular states (e.g., edge currents) are important for (e) both the 'integer' and 'fractional' effects. |
| (317) In addition, Coulomb interaction is also essential in (eb) the fractional quantum Hall effect. |
| (318)The observed strong similarity between integer and (=) fractional quantum Hall effects is explained by (e) the tendency of electrons to form bound states with (e&eb) an even number of magnetic flux quanta, called composite fermions. |
| (319)Duality is the circumstance in which cranking up interactions results in (eb) another system we understand, usually one that can be described in terms of (e&eb) another set of weakly coupled particles. |
| (320)When such a dual representation is found, (eb) it gives a solution to the strong coupling problem. |
| (321) A further reason that theoretical physicists get excited about dualities is that they undermine (e) the very notion of an “elementary particle.” |
| (322)In terms of the original variables, the dual, weakly coupled particles are (=) generally solitons—large, heavy, classical objects made up of (e) many quanta. |

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| (323)The remarkable thing is that solitons can become (=) "elementary," in the sense that the model that describes the system can be reformulated with the solitons as (=) the basic constituents. |
| (324)An example of a duality, which we will use below, arises in (eb) Maxwell's theory of electromagnetism, where (e) the strength of the interactions is determined by (e) the electric charge e. |
| (325)The duality exchanges electric and (e&eb) magnetic point charges, and the interaction strength in the dual description is 1⁄e, which is small when (e) e is large. |
| (326)This idea has been generalized to (e) many other field theories, thanks, in particular, to tools of (e) supersymmetry [4]. |
| (327)Another classic example of duality involves (e&eb) vortices. |
| (328) A vortex is a localized object in the plane around which a phase field winds (e&eb). (A phase field takes values on the circle.) |
| (329)Dualities relating particles and vortices have (e) a rich history in particle physics and condensed-matter physics. |
| (330) The simplest version relates the particles in the normal state of a superfluid to (e&eb) the vortices in a dual superconductor [5]. |
| (331)This charge-vortex duality has been (=) a workhorse in the study of two-dimensional condensed-matter systems, such as (=) fractional Quantum Hall States, superconducting thin films, and beyond-Landau critical phenomena [6]. |
| (332)Vortices can be (=) fermions (as they frequently are, for instance, in type-II superconductors). |
| (333) A dual description of a system with (e&eb) fermionic vortices thus requires (e) a new ingredient, since (e) it must include (e) such a fermionic particle. |
| (334)The new work [1–3] builds on (e&eb) earlier theoretical work [7] to make (eb) an appealing proposal for (e) the required duality: the fermionic vortex of QED in (eb) 2+1 dimensions is (=) a free Dirac fermion. The argument for this proposal builds on our understanding of TIs. |
| (335)TIs are distinguished only by (e) their edge states. |
| (336) The most famous example is a (3+1)-dimensional time-reversal invariant TI, which can host boundary states whose band structure is defined by (e) a single Dirac cone. |
| (337)A single time-reversal-symmetric Dirac cone is not possible without (e) the bulk TI, so its appearance in experiments provides (eb) a signature of the bulk phase. |
| (338) But the Dirac cone is not (e) the only possible signature. |
| (339) Deformations of the boundary conditions can dramatically change (e&eb) the spectrum of the edge states, yet preserve some essential "TI-ness." This essential property is (=) called an anomaly in high-energy theory. |
| (340)Max Metlitski of the University of California, Santa Barbara, and Ashvin Vishwanath of the University of California, Berkeley [1], and independently, Chong Wang and Todadri Senthil at the Massachusetts Institute of Technology, Cambridge [2], use this ambiguity in the edge theory and the electric-magnetic duality of an auxiliary theory of bulk electromagnetism coupled to (e&eb) the TI to argue for the duality. |
| (341)When I said that (3+1)-dimensional QED had a duality that interchanged the electric and magnetic charges, one could have (e) complained that the electron is (=) a fermion, while (e) a magnetic monopole may not be. |
| (342)However, the authors observe that in a TI, a charge-two monopole is (=) also a fermion. |

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| (343) By analyzing the spectrum of electric and magnetic charges, they conclude that the duality exchanges the electron and (e&eb) the charge-two monopole and, moreover, exchanges a TI with (e&eb) time-reversal symmetry with (e&eb) a "chiral TI," which instead has (e) particle-hole symmetry |
| (344). But this is just what is required of (e) a symmetry that exchanges a fermionic particle with (e&eb) a fermionic vortex: time-reversal symmetry maps the electron to itself but maps the vortex to (e&eb) another vortex whose phase winds (e&eb) the other way (see Fig. 1). |
| (345)The exchange of particles and holes, in turn, preserves (eb) the vortex. |
| (346)The bulk duality implies (eb) a duality between edge theories. |
| (347)The freedom to deform the edge of the TI implies (eb) a weak form of equivalence: although the edge theories are (=) topologically the same (in terms of symmetry assignments and (e&eb) anomalies), they need not have (e) identical low-energy physics. |
| (348)One side of the duality is (=) a free, massless Dirac particle, but the long-distance physics of (2+1)-dimensional QED is (=) a long-standing problem, which is important in the study of spin liquids. |
| (349) Many gauge theories in two dimensions develop (eb) a gap in their energy spectrum because of (e) the proliferation of monopoles. |
| (350) In the realization proposed by Metlitski and Vishwanath [1] and by Wang and Senthil [2], this cannot happen because (e) the monopoles carry electric charge, and indeed the stronger notion of duality as (=) an equivalence of low-energy physics may hold. |
| (351)Evidence for the stronger relation comes from (e) David Mross and colleagues at the California Institute of Technology, Pasadena [3], who have derived a lattice version of the duality using (e) a clever but explicit change of variables. |
| (352)Their construction is particularly interesting because (e) it uses (e) a lattice of one-dimensional models which themselves can only occur at the edge of a TI. |
| (353)Authors explicitly derive the duality between a free electronic Dirac cone and (e&eb) quantum electrodynamics in (eb) (2+1) dimensions (QED3) with (e&eb) N=1 fermion flavors. |
| (354)The duality proceeds via (e&eb) an exact, nonlocal mapping from electrons to (e&eb) dual fermions with (e&eb) long-range interactions encoded by (e) an emergent gauge field. |
| (355)This mapping allows (eb) us to construct (eb) parent Hamiltonians for exotic topological-insulator surface phases, derive (eb) the particle-hole-symmetric field theory of (e) a half-filled Landau level, and nontrivially constrain (e) QED3 scaling dimensions. |
| (356)They similarly establish duality between bosonic topological insulator surfaces and (e&eb) N=2 QED3. |

**Black hole thermodynamics with generalized uncertainty principle Li Xiang, X. Q. Wen (Submitted on 6 Jan 2009 (v1), last revised 18 Mar 2009 (this version, v2))**

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| (357) In the standard viewpoint, the temperature of a stationary black hole is proportional to (=) its surface gravity, $T_H=\hbar\kappa/2\pi$. |
| (358)This is a semiclassical result and the quantum gravity effects are not (e) taken into consideration. |
| (359)This Letter explores a unified expression for (e) the black hole temperature in the sense of (e) a generalized uncertainty principle (GUP). |
| (360) Discussion involves a heuristic analysis of (e) a particle which is absorbed by (e) the black hole. |
| (361)Besides a class of static and spherically symmetric black holes, an axially symmetric Kerr-Newman black hole is considered. Different from the existing literature, authors suggest (eb) that the black hole's irreducible mass represent (eb) the characteristic size in (eb) the absorption process. |

(362)The information capacity of a remnant is also discussed by (e) Bousso's D-bound in de Sitter spacetime. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP 0910:046,2009 DOI: 10.1088/1126-6708/2009/10/046 Cite as: arXiv:0901.0603 [gr-qc] (or arXiv:0901.0603v2 [gr-qc] for this version)

IF we define

$$[a_7]_7 = [f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}]\,P_7^*$$

$$[a_7]_7 = \left[f_{7*}^{7}\right] P_7^* - \frac{h_*^7}{h_*} P^*(0)$$

$$e = [g_*^7]_7\, Q_7^*$$

$$\eta_7(t) = \left[h_*\right]_7 \beta_7 \left(\theta_7(t)\right) - \mu_7(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\left[a_7\right]M_7 - [a_7]_7 W_7 - [n(t)]M_7$$

$$w_7(t) = eM_7$$

Write

$$A = \begin{pmatrix} -[a_7]_7 - [a_7]_7 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)= \begin{pmatrix} [\eta(t)]_7 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_7 \\ W_7 \end{pmatrix}$$

The solution of (32)is

$$X_7(t) = e^{A,t}[X]_7(0) + \int_0 e^{At}$$

## PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION

## ATTRITION MODELS

## MODULE EIGHT

Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal.

### 25. MODEL FORMULATION

**Model is applied to systems given in the following:**

**Researchers blur the line between classical and quantum physics by connecting chaos and entanglement July 12, 2016 by Sonia Fernandez**

(1557) Experimental link between quantum entanglement (left) and (e&eb) classical chaos (right) found using (e) a small quantum computer. Credit: University of California - Santa Barbara Using a small

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| quantum system consisting of three superconducting qubits, researchers at UC Santa Barbara and Google have uncovered a link between aspects of classical and (e&eb) quantum physics thought to be unrelated: classical chaos and (e&eb) quantum entanglement. |
| (1558) Their findings suggest (eb) that it would be possible to use (e) controllable quantum systems to investigate certain fundamental aspects of nature. |
| (1559) "It's kind of surprising because (e) chaos is this totally classical concept—there's (=) no idea of chaos in a quantum system," Charles Neill, a researcher in the UCSB Department of Physics and lead author of a paper that appears in Nature Physics. |
| (1560) "Similarly, there's no concept of entanglement within (eb) classical systems. And yet it turns out that chaos and (e&eb) entanglement are really very strongly and clearly related." |
| (1561) Initiated in the 15th century, classical physics generally examines and describes (eb) systems larger than atoms and molecules. |
| (1562) It consists of hundreds of years' worth of study including Newton's laws of motion, electrodynamics, relativity, thermodynamics as well as chaos theory—the field that studies the behavior of (e) highly sensitive and unpredictable systems. |
| (1563) One classic example of chaos theory is the weather, in which a relatively small change in one part of (e) the system is enough to foil predictions—and vacation plans—anywhere on the globe. |
| (1564) At smaller size and length scales in nature, however, such as those involving (e&eb) atoms and photons and their behaviors, classical physics falls short. |
| (1565) In the early 20th century quantum physics emerged, with its seemingly counterintuitive and sometimes controversial science, including (e) the notions of superposition (the theory that a particle can be located in (eb) several places at once) and entanglement (particles that are deeply linked behave as (=) such despite physical distance from one another). |
| (1566) And so began the continuing search for connections between the two fields. All systems are fundamentally quantum systems, according Neill, but the means of describing in a quantum sense the chaotic behavior of, (e) say, air molecules in an evacuated room, remains limited. |
| (1567) Imagine taking a balloon full of air molecules, somehow tagging them so you could see them and then releasing them (eb) into a room with no air molecules, noted co-author and UCSB/Google researcher Pedram Roushan. |
| (1568) One possible outcome is that the air molecules remain clumped together (e&eb) in a little cloud following the same trajectory around the room |
| (1569) And yet, he continued, as we can probably intuit, the molecules will more likely take off in (eb) a variety of velocities and directions, bouncing off walls and interacting with (e&eb) each other, resting after the room is (=) sufficiently saturated with them. |
| (1570) "The underlying physics is (=) chaos, essentially," he said. |
| (1571) The molecules coming to rest—at least on (eb) the macroscopic level—is the result of (e) thermalization, or of reaching equilibrium after they have achieved (eb) uniform saturation within the system. |
| (1572) But in the infinitesimal world of quantum physics, there is (=) still little to describe that behavior. |
| (1573) The mathematics of quantum mechanics, Roushan said, do not allow for (e) the chaos described by Newtonian laws of motion. |
| (1574) To investigate, the researchers devised an experiment using (e) three quantum bits, the basic computational units of (e) the quantum computer. |

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| (1575) Unlike classical computer bits, which utilize (e) a binary system of two possible states (e.g., zero/one), a qubit can also use (e) a superposition of both states (zero and one) as a single state. |
| (1576) Additionally, multiple qubits can entangle, or link so closely that their measurements will (=) automatically correlate. |
| (1577) By manipulating these qubits with electronic pulses, Neill caused them to interact, rotate and evolve (e&eb, e&eb, eb) in the quantum analog of (e) a highly sensitive classical system. |
| (1578) The result is a map of entanglement entropy of (e) a qubit that, over time, comes to strongly resemble that of (e) classical dynamics—the regions of entanglement in (eb) the quantum map resemble the regions of (e) chaos on the classical map. |
| (1579) The islands of low entanglement in (eb) the quantum map are located in the places of (e) low chaos on the classical map. |
| (1580) "There's a very clear connection between entanglement and (e&eb) chaos in these two pictures," said Neill. |
| (1581) "And, it turns out that thermalization is (=) the thing that connects chaos and (e&eb) entanglement. |
| (1582) It turns out that they are actually the driving forces behind (eb) thermalization. |
| (1583) "What we realize is that in almost any quantum system, including on quantum computers, if you just let it evolve and you start to study what happens as a function of time, (eb) it's going to thermalize," added Neill, referring to the quantum-level equilibration. "And this really ties together the intuition between classical thermalization and (e&eb) chaos and how it occurs (eb) in quantum systems that entangle." |
| (1584) The study's findings have fundamental implications for (e) quantum computing. |
| (1585) At the level of three qubits, the computation is relatively simple, said Roushan, but as researchers push to build increasingly sophisticated and powerful quantum computers that incorporate (e) more qubits to study highly complex problems that are beyond the ability of classical computing—such as those in the realms of machine learning, artificial intelligence, fluid dynamics or chemistry—a quantum processor optimized for such calculations will be a very powerful tool. |
| (1586) "It means we can study things that are completely impossible to study right now, once we get to bigger systems," said Neill.More information: Nature Physics, DOI: 10.1038/nphys3830 , http://www.nature.com/nphys/journal/vaop/ncurrent/full/nphys3830.html Journal reference: Nature Physics search and more info website Provided by: University of California - Santa Barbara |

**Title: coherent control of macroscopic quantum states in a single-Cooper-pair box Authors: Nakamura, Y.; Pashkin, Yu. A.; Tsai, J. S. Publication: Nature, Volume 398, Issue 6730, pp. 786-788 (1999). (Nature Homepage) Publication Date: 04/1999 Origin: NATURE DOI: 10.1038/19718 Bibliographic Code: 1999Natur.398...786N**

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| (1587) A nanometre-scale superconducting electrode connected to (e&eb) a reservoir via (e&eb) a Josephson junction constitutes (e) an artificial two-level electronic system: a single-Cooper-pair box. |
| (1588) The two levels consist of (e) charge states (differing by 2e, where (e) e is the electronic charge) that are coupled by (e) tunnelling of Cooper pairs through (e&eb) the junction. |
| (1589) Although the two-level system is macroscopic, containing (e) a large number of electrons, the two charge states can be (=) coherently superposed. |

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| (1590) | The Cooper-pair box has therefore been suggested as (=) a candidate for a quantum bit or `qubit'- the basic component of a quantum computer. |
| (1591) | Here authors report the observation of quantum oscillations in (eb) a single-Cooper-pair box. |
| (1592) | By applying a short voltage pulse via (e&eb) a gate electrode, we can control (e&eb) the coherent quantum state evolution: the pulse modifies (e&eb) the energies of (e) the two charge states non-adiabatically, bringing them into (e&eb) resonance. |
| (1593) | The resulting state-a superposition of (e) the two charge states-is detected by (e) a tunnelling current through (e&eb) a probe junction. |
| (1594) | Results demonstrate electrical coherent control of (e) a qubit in a solid-state electronic device. |

**Superconducting quantum computing From Wikipedia, the free encyclopedia**

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| (1595) | Superconducting quantum computing is a promising implementation of (e) quantum information technology that involves (e&eb) nanofabricated superconducting electrodes coupled through (e&eb) Josephson junctions. |
| (1596) | As in a superconducting electrode, the phase and (e&eb) the charge are conjugate variables. |
| (1597) | There exist three families of superconducting qubits, depending on (=) whether the charge, (e&eb) the phase, or neither of (e) the two is good quantum number. |
| (1598) | These are respectively termed charge qubits (e&eb) flux qubits, and (e&eb) hybrid qubits. |

**Theory**

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| (1599) | Unlike many other physical implementations of a qubit which involve (e&eb) exclusively two-level systems (such as nuclear spin and (e&eb) photon polarization), the integrated quantum circuit involved in (e&eb) a superconducting qubit is (=) a multi-level system, of which only the first two levels are used as (=) the computational basis. |
| (1600) | A basic requirement for such an implementation is that the energy levels are (=) not uniformly spaced, so that (e) photons of a particular frequency which cause (eb) transition between the 0 and 1 levels do not(e) cause transitions from the first level to (e&eb) the higher levels as well. |
| (1601) | For electronic signals to be carried from one part of the circuit to (e&eb) another without energy loss and (e&eb) hence decoherence, the system also needs to be (=) non-dissipative, i.e., the metallic parts involved should have (e) zero resistance. |
| (1602) | These circuits currently need to be (=) operated at very low temperatures so that, 1) superconductivity is (=) realized, and 2) thermal fluctuations do not (e) cause transitions between (e&eb) energy levels. |

**Quantum LC Circuit**

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| (1603) | The simplest non-dissipative quantum circuit consists simply of (e) an inductor L {\displaystyle L} L and (e&eb) capacitor C {\displaystyle C} C, with (e&eb) the metallic wires connecting them being superconducting. |
| (1604) | Here the flux Φ {\displaystyle \Phi } \Phi through (e&eb) the inductor and (e&eb) the charge Q {\displaystyle Q} Q in the capacitor are (=) canonically conjugate variables which obey (eb) the commutation relation [ Φ , Q ] = i ℏ {\displaystyle [\Phi ,Q]=i\hbar } [\Phi ,Q]=i\hbar |

**The Hamiltonian associated with this system is**

(1605) H = Φ 2 2 L + Q 2 2 C {\displaystyle H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}} H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}

(1606) This gives rise to the energy levels E = ℏ ω 0 ( n + 1 2 ) {\displaystyle E=\hbar \omega _{0}(n+{\frac {1}{2}})} E=\hbar \omega _{0}(n+{\frac {1}{2}}) Here ω 0 = 1 L C {\displaystyle \omega _{0}={\frac {1}{\sqrt {LC}}}} \omega _{0}={\frac {1}{{\sqrt {LC}}}}

**Quantum gravity Carlo Rovelli (2008), Scholarpedia, 3(5):7117. doi:10.4249/scholarpedia.7117 revision #135539 [link to/cite this article] Post-publication activity Curator: Carlo Rovelli Dr. Carlo Rovelli, Center of Theoretical Physics, and Marseille, France**

(1607) Quantum Gravity is the name given to any theory that describes gravity in (eb) the regimes where quantum effects cannot be disregarded.

(1608) At present, there is no such a theory which is universally accepted and confirmed by experience. Therefore the term "Quantum Gravity" indicates (eb) more an open problem than a specific theory.

(1609) Several research lines, at different levels of development, offer tentative solutions to the problem. These tentative quantum-gravity theories are variously viewed as competing research directions, or as contributions to the common goal of finding the physically correct theory. The quest for the good quantum theory of gravity bears on a number of fundamental issues and it is sometimes presented as the most important open problem in fundamental physics: the "Holy Grail" of contemporary theoretical physics.

**The problem of quantum gravity**

(1610) The current theory of (e) gravity is general relativity, a theory developed by Albert Einstein at the beginning of the XX century, which has changed in depth our understanding of (e) what are space and time.

(1611) Similarly, quantum mechanics has changed in depth our understanding of matter, energy and causality. These two theories, today very well empirically supported, have opened (eb) a major conceptual revolution in physics.

(1612) However, they appear to be incompatible, at least at a first reading, because (e) each of the two is formulated on the basis of (e) principles that are explicitly contradicted by (e) the other theory.

(1613) Therefore in spite of the immense development of scientific knowledge it has delivered, the physics of the XX century has left us with (e&eb) a major confusion on what are the basic conceptual ingredients for (e) understanding the physical world.

(1614) In other words, the scientific revolution opened by general relativity and quantum mechanics at the beginning of (e) the XX century is not concluded yet, and a new synthesis is required. Quantum gravity, merging general relativity and quantum mechanics should be this new synthesis.

**Quantum space and quantum time**

(1615) Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.

(1616) In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.

(1617) They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.

(1618) When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.

(1619) Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.

**The Planck scale**

(1620) Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".

(1621) Here $\hbar$ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.

(1622) The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.

(1623) Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)

(1624) Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.

## Physical applications

Some phenomena where quantum gravity plays a major role are the following.

**The microscopic structure of spacetime**

(1625) If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.

(1626) Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length

(1627) Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.

(1628) A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.

**Early cosmology**

(1629) According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.

(1630) However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.

(1631) A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.

(1632) Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".

(1633) One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.

**Black holes**

(1634) Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.

(1635) First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.

(1636) Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.

(1637) What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.

(1638) Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.

**Astrophysical effects**

(1639) Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.

(1640) An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.

(1641) The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.

**Main approaches**

(1642) The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.

(1643) Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.

(1644) The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.

**String theory**

(1645) String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.

(1646) Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.

(1647) String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.

(1648) Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.

(1649) However, when summed the entire series appears to be (=) divergent.

(1650) In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.

(1651) Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.

(1652) Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.

(1653) In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.

**Loop quantum gravity and spin foams**

(1654) Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.

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| (1655) The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity. |
| (1656) Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime. |
| (1657) The granular properties of space can be (=) explicitly computed. |
| (1658) In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory. |
| (1659) The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks. |
| (1660) Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure. |
| (1661) The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit. |
| (1662) Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism. |
| (1663) This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism. |

**Noncommutative geometry**

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| (1664) Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime. |
| (1665) Quantum mechanics teaches us that dynamical quantities are (=) noncommutative --meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed. |
| (1666) It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry. |
| (1667) A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces. |
| (1668) Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces. |
| (1669) This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity. |

**Others**

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| (1670) Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays |

(1671) . Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection.

(1672) Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible.

(1673) The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.

Let $p_8(t, a_8)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(g_8)Q_8$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_8(0) = 0, \quad g_8(P_8) > 0, \quad P_8 \; 0, \qquad \lim_{P_8 \to \infty} g_8(P_8) = (K_8)_8 \; \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_8 + (\rho_a)_8 + D_8(a_8, P_8, Q_8)\rho_8 + 0, \qquad 0 < a_8 \; \infty,$$

$$\rho_8(t,0) = \int_0^\infty n_8(a_8, P_8, Q_8)\rho_8(t_8, a_8)da_8$$

Where $D_8$ is the mortality function and $n_8$ is the fertility function.

Put $_8$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_8}{dt}\right]_8 = x_8\left[\frac{dQ_8}{dt}\right]_8 - x_8\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_8}{dt}\right]_8 = x_3\left[\frac{dQ_3}{dt}\right]_8 - x_4\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_8}{dt}\right]_8 - x_6\left[\frac{dQ_3}{dt}\right]_8$$

Where $x_8, x_8, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_8$ and $g_8(P_8)Q_8$

Now when you add the term $(b_8Q_8)$ with $g_8(P_8)(Q_8)$ we are infect see is that $(b_8Q_8)$ is credited incrementally part of $g_8(P_8)(Q_8)$ or instant nearly, when subtracted from $\frac{dQ_8}{dt}$ , it means that $g_8(P_8)(Q_8)$ is slowly eaten by $\frac{dQ_8}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_8 = \mu_8(a_8) + (f_8)_8(P_8)_8 + (f_8)_8(Q)_8 \quad (2)$$

$$n_8 = \beta_8(a_8)h_8(Q_8), \quad (3)$$

Where $\mu_8$the age- dependent debit rate whit is $\mu_8(a_8)$ 0 $\forall a > 0$; $(f_8)_8$ is the density- dependent death rate with $f_8(0) = 0$, $(f_8^8)_8(P_8)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_8)_8(P_8) = (k_8)_8$ ; $f_8$ is the functional response of creditors to prey with $[f_8]_8(0) = 0, (f_8^8)_8(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_8)_8(Q_8) = [k_8]$ ; $\beta$ is the age

dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^8(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_8 + (\rho_a)_8 = -\mu_8(a_8) + (f_8)_8 P_8) + \big(f_8)_8(Q_8)\big)(\rho_8)_8 \quad (4)$$

$$\rho_8(t_8, 0) = h_8(Q_8) \int_0^\infty \beta_8 (a_8)\rho_8(t, a_8)da, \quad (5)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(P_8)Q_8 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_8$ and $Q_8$ of (4-6) satisfy

$$0 \ \rho_8(t, a_8) \ \rho_8^*(t, a_8),$$

$$0 \ Q_8(t) \ Q_8^*(t),$$

Where $\rho_8^*(t, a_8)$ and $Q_8^*(t)$, are solutions of

$$(\rho_t^*)_8 + (\rho_a^*)_8 = -(\mu_8(a_8) + (f_8)_8(P_8^*) + (f_8)_8 Q_8^*))\ \rho_8^*, \quad (7)$$

$$\rho_8^*(t, 0) = \int_0^\infty \beta_8(a_8)\ \rho_8^*(t, a_8)da_8 \quad (8)$$

$$\frac{dQ_8^*}{dt} = -b_8\ Q_8^* + g_8(P_8^*)Q_8^*, \quad (9)$$

With $P_8^*(t) = \int_0^\infty \rho_8^*\ (t, a_8)da_8$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_8^*(t)}{dt} = P_8^*(t)\big(r_8 - (f_8)(P_8^*) - (f_8)_8(Q^*)\big), \tag{10}$$

$$\frac{dQ_8^*(t)}{dt} = Q_8^*(t)(-b_8 + g_8(P_8^*)), \tag{11}$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_8(a)e^{-\int_0^a \mu_8(\tau)dt}e^{-ra}\ da_8 = 1. \tag{12}$$

If $r < 0$, the trivial solution $P_8^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(t)dt}\ da_8 > 1, \tag{$H1$}$$

Which guarantees $r_8 > 0$.

*Remark.* If $\beta_8(a) = [\beta_0]_8 a_8 e^{-\alpha_8 a_8}, \mu_8(a_8) \equiv \mu_8$ then the solution of (12) is

$$r_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_8]_8(P_8) = 0, [f_8](Q_8) = \lambda_8 Q_8\ h_8(Q_8) = 1, \qquad and\ \ g_8(P_8) = c_8 P_8$$

$$[\rho_8]_8 + [\rho_a]_8 = -(\mu_8 + \lambda_8 Q_8)\rho_8$$

$$\rho_8(t,0) = \int_0^{\infty} [\beta_0]_8 a_8 e^{-\alpha_8 a_8} \rho_8(t,a_8)\, da_8$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + c_8 P_8 Q_8$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_8}{dt} = \omega_8 P_8 - \lambda_8 P_8 Q_8$$

$$\frac{dQ_8}{dt} = b_8 Q_8 - c_8 P_8 Q_8$$

Where $\omega_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$

## 26. THE EQUILIBRIUM $(\rho_0)_8(a_8), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_8(a_8), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_8 + (\rho_a)_8 = -\big(\mu_8(a_8) + (f_8)_8(P_8)\big)\rho,$$

$$\rho_8(t,0) = \int_0^{\infty} \beta_8(a_8)\rho_8(t,a_8)da_8 \qquad (13)$$

The solution of (13) has the form of

$$\rho_8(t,a_8) = Ce_8^{r(t-a_8)} - \int_0^a (\tau)d\tau - \int_0^8 (f_8)_8\big(P_8(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_8(t) = C_8 e^{r_8 t - \int_0^t (f_8)_8\left(P_8(\tau)\right)d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau}\, da.$$

Hence, there exists an equilibrium $P_0 = \int_8^{-8}(r)$ if $k_8 > r$. This is equivalent to

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{(k_8)_8 a_8} da_8 < 1.$$

In fact, if we define

$$R_8(s_8) = \int_0^\infty \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau)d\tau} e^{-s_8 a_8}\, da_8$$

$R(s)$ is a decreasing function. $R_8(r_8) = 1$ And $R_8(k_8)_8 < 1$ imply $r_8 < (k_8)_8$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_8(a_8), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{k_8 a_8} da_8 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_8 \equiv \int_0^\infty (\rho_0)_{88}(a_8)da_8$ *satisfies*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{-(f_8)_8(P_0)a_8} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_8(a_8), 0)$

Let $u_8(t,a_8) = \rho(t,a_8) - \rho_0(a_8)$ Then

$$N_8(t) \equiv \int_0^\infty u_8(t,a_8)da_8 = \int_0^\infty \rho_8(t,a_8)da_8 - \int_0^\infty (\rho_0)_8(a_8)da = P_8(t) - P_0)_8$$

And

$$(u_8)_8 + (u_a)_8 = (\rho_t)_8 + (\rho_a)_8 - (\rho_0^8)_8(a_8)$$

$$= -\mu_8(a_8) + (f_8)_8(P_8) + (f_8)_8(Q_8))\,(u_8 + (\rho_0)_8) + \mu_8(a_8)(\rho_0)_8 + (f_8)_8(P_0)_8(\rho_0)_8$$

$$= -\mu_8(a_8)u_8 - (f_8)_8(P_8) + (f_8)_8(Q_8))\,u_8 - (\rho_0)_8\,(f_8)_8(P_8) - (f_8)_8(P_0)_8$$

$$+ (f_8)_8(Q_8)), \qquad (15)$$

Where $(\rho_0^8)_8(a_8)$ denotes $d(\rho_0)_8(a_8)/da_8$

Expansion $(f_8)_8(P_8)$ and $(f_8)_8(Q_8)$ at $(P_0)_8$ and 0 yield respectively

$$(f_8)_8(P_8) = (f_8)_8((P_0)_8 + N_8)$$

$$= (f_8)_8(P_0)_8 + (f_8^8)_8(P_0)_8N_8 + o_8(N_8), \qquad (16)$$

$$(f_8)_8(Q_8)$$

$$= (f_8^8)_8(0)Q + o_8(Q_8) \qquad (17)$$

Write $f_8(P_0) = f_{80}, f_8'(P_0) = f_{80}'$, and $f_8'(0) = f_{80}'$ and substitute (16) and (17) into (15) to obtain

$$(u_8)_8 + (u_a)_8 \approx -(\mu_8(a_8) + (f_{80})_8)u - (f_{80}^8)_8(\rho_0)_8(N_8) - (f_{80}^8)_8(\rho_0)_8(Q_8)$$

$$u_8(t,0) = \rho_8(t,0)(\rho_0)_8(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a) + (\rho_0)_8(a_8)da_8 - (\rho_0)_8\,(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a_8)da_8 + (h_8(Q_8) - 1)(\rho_0)_8\,(0) \qquad (18)$$

Expanding

$$h_8(Q_8) = 1 + (h^8)_8(0)Q_8 + 0(Q_8) \equiv 1 + h_0^{'}Q_8 + 0Q_8$$

And substituting it into (18), we arrive at

$$u_8(t,0) \approx \int_0^{\infty} \beta_8\,(a_8)u_8(t,a_8)da_8 + (h_{a_8}^{'})(\rho_0)_8 Q_8$$

The linearization is the following :

$$[u_8] + [u_8] = -(\mu_8(a_8) + [t_{80}]_8 u_8 - [\alpha_8]_8(a_8)N_8 - [\alpha_8]_8(a_8)Q_8$$

$$u_8(t,0) = \int_0^{\infty} \beta_8(a_8)u_8(t,a_8)da_8 + [\xi_8]Q_8 \qquad (19)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8$$

where

$$[\alpha_8]_8(a_8) = [t_8^8 a]_8[p,o]_8(a_8)d2(a_8) = [f_{80}^8]_8(\rho\, o\,)_8(a_8)\xi = [\mathrm{h\ o}^8]_8[\rho \mathrm{o}]_8(\mathrm{o}) \text{ and } \mathrm{b}_8 = \mathrm{b}_8 - g_8(P_8)$$

If $g_8(P_0) > b$, clearly, $(\rho_0)_8(a_8), 0)$ is unstable and the predator population is persistent. Assume $g_8(P_0)_8 < b$.

Then, $Q_8(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_8 + u_a = -(\mu_8(a_8) + (f_{80})_8)u_8 - (\alpha_8)_8(a_8)N_8$$

$$= -(\mu_8(a_8) + f_{80})u_8 - \alpha_8(a_8)\int_0^{\infty} u_8(t,a_8)da_8$$

$$u_8(t,0) = \int_0^\infty \beta_8\,(a_8)u_8(t,a_8)da_8$$

Using the method developed by Cushing [2, 3], we suppose

$$u_8(t,a_8) = \tilde{u}_8(a_8)^{e^{\gamma(t-a_8)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_8}\,da_8$. Then,

A solution of the first differential equation for $\tilde{u}(a_8)$ gives,

$$\tilde{u}_8(a_8) = e^{-\int_0^a \mu_8(\tau_8)d\tau - (f_{80})_8 a}\left(\tilde{u}_8(0) - w\int_0^a e^{\int_0^\tau \mu_8(s_8)ds + (f_{80})}\tau_8\,(\alpha_8)_8(\tau_8)e^{\gamma_8\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\right)\tilde{u}_8(0)$$

$$+w\int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\int_0^a (\alpha_8)_8(\tau_8)e^{-\gamma_8\tau_8}e^{(f_{80})_8\tau_8 + \int_0^\tau \mu_8(s_8)ds}d\tau_8\,da_8$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}da\right)(\tilde{u}_8(0))_8$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_8 a_8-(f_{80})_8 a_8-\int_0^a \mu_8(\tau_8)d\tau}\int_0^a e^{(f_{80})_8\tau_8+\int_0^\tau \mu_8(s_8)ds+y\tau}\alpha_8(\tau)d\tau\, da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_8 a_8-(f_{80})_8 a_8-\int_0^a \mu_8(t_8)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_8\,(a_8)Y_8(\gamma_8,a_8)da_8\right)u_8(0)$$

$$+w\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)\int_0^a(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8=0,$$

(24)

$$\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\tilde{u}_8(0)$$

$$+\left(-1-\int_0^{\infty}Y_8(\gamma_8,a_8)\int_0^a(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8\right)w=0.$$

Equation (20) has a solution of the form of $u_8(t,a_8)=u_8(a_8)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_8(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_8(\gamma_8,a_8)=\int_0^a(\rho_0)_8(\tau_8)Y^{-8}\,(\gamma_8,a_8)d\tau_8$. We obtain the characteristic equation

$$f_{80_8}^{l}\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\left(\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$+\left(1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)da\right)\left(1+(f_{80}^{l})_8\int_0^{\infty}Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_8(a_8) = (\rho_0)_8(0)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - (f_{80})_8 a_8}$$

$$I_8(\gamma_8, a_8) = \int_0^a (\rho_0)_8(\tau_8)e^{\gamma_8\tau_8 + (f_{80})_8\tau + \int_0^\tau \mu_8(s_8)ds_8}d\tau$$

$$= \frac{(\rho_0)_8(0)}{\gamma_8}(e^{\gamma_8 a_8} - 1)$$

Let

$$E_8(a_8) = e^{-(f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau_8}$$

Then $[Y_8]_8[\gamma_8]_8, a_8) = E_8(a_8)e^{-\gamma_8 a_8}$. by a simple calculation,

$$\int_0^\infty \beta_8(a_8)Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma}\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}\ da\right);$$

$$\int_0^\infty Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma_8}\left(\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}\ da_8\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_8(0)}{\gamma_8}[f_{80}^8]\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}da_8\ (\int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8)$$

$$+\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right).$$

$$\left(1+\frac{(\rho_0)_8(0)}{\gamma_8}(f_{80}^{'})_8\left(\int_0^{\infty}E_8(a_8)da_8-\int_0^{\infty}E_8(a_8)e^{-\gamma_8 a_8}da_8\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right)(1+\frac{(f_{80}^{'})_8}{\gamma_8}(\rho_0)_8(0)\int_0^{\infty}E_8(a_8)da_8)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)da_8=1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\neq 0,$$

Since $\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)\int_0^{a}\alpha_8(\tau_8)Y^{-8}(\gamma_8,\tau_8)d\tau_8 da_8>0$. Therefore, the characteristic solution is

$$\gamma_8=-(f_{80}^{'})_8(\rho_0)_8\int_0^{\infty}E_8(a_8)da_8.$$

Since $f_8^{'}(P_0)_8>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_8),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_8),0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_8 + (\mu_8(a_8) + [f_8]_8(P_8^*) + [f_8]_8(Q_8^*))P_8^* = 0, \quad (26)$$

$$p_8^*(0) = h_8(Q_8^*)\int_0^\infty \beta_8(a_8)p_8^*(a_J)p_8^*(a_8)da_8 \quad (27)$$

$$g_8(P_8^*) - b_8 = 0, \quad (28)$$

Where $P_8^* = \int_0^\infty p_8^*(a_8)\, da_8$

Equation (28) yields $P^* = g^{-8}(b)$,

From (26),

$$p_8^*(a) = p_8^*(0)e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}$$

Substituting into (27), we need to solve

$$\left(1 - h_8(Q_8^*)\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}da_8\right)p_8^*(0) = 0$$

For $Q_8^*$

Fix $p_8^*(0)$ and $P_8^*$ Let

$$H_8(z_8) = 1 - h_8(z_8)\int_0^\infty \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}da_{,8}$$

And assume that

$$[k_3]_8 \int_0^{\infty} \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(g_8^{-8}(b_8))a_8 - [k_8]_8 a_8}\, da_8 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty} [H_8(z_8)] > 0.$$

Clearly,

$$\frac{dH_8}{dz_8} = h_8(z)(\,[f_8^{'}]_8(z_8)\,) \int_0^{\infty} a_8 \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - f_8(P_8^*)a_8 - f_8(z_8)a_8} da_8$$

$-h^{'}(z_8) \int_0^{\infty} \beta_8(a_8)\, e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}\; da_8$

By the assumptions on $\left[f_8^8\right]_8 and\ [h^8]^8 dH_8/dz > 0$. Hence, if

$$H_8(0) = 1 - \int_0^{\infty} \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8} da_8 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_8(0) < 0$ follows if $g_8^-(b_8) < [f_8^{-8}]_8(r_8)$, where $r_8$ is the solution of (12).

Indeed, since

$$R_8(s_8) = \int_0^{\infty} \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 e^{-y_8^a}} da_8$$

Is a decreasing function of $s_8$ when $P_8^* = [g^{-8}]_8(b_8) < [f_8^{-8}]_8(r_8)$, that is $[f_8]_8(P_8^*) < r$, $R\big(f_8(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_8(g^{-8}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8\left(g_8^{-8}(b)\right)a_8} > 1. \qquad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_8^*(a_8), Q_8^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_8^*(a_8), Q_8^*)$ is more complicated. We again consider perturbations $m_8(t_8 a_8) = p_8(t_8, a_8) - p_8^*(a_8)$ and $W_8(t) = Q_8(t) - Q^*$.Then

$$M_8(t) \equiv \int_0^\infty m_8(t, a_8)da_8 = P_8(t) - P_8^*,$$

$$[m_8]_8 + [m_a]_8 = [p_8]_8 + [p_a]_8 - \left[p^*\right]_8(a_8) \approx -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8$$

$$-[f_{8*}^8]_8 p_8^*(a_8)M - [f_{8*}^8]_8 p_8^*(\mathrm{a})W_8,$$

$$m_8(t,0)p_8(t,0) - p_8^*(0) \approx h_* \int_0^\infty \beta_8(a_8)m_8(t,a_8)\,da + \frac{[h_*^8]_8}{[h_*]_8}\, p_8^*(0)W_8$$

$$\frac{dW_8}{dt} = \frac{dQ_8}{dt} \approx [g_*^8]_8 Q_8^* M_{,8}$$

Where $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $\left[f_{8*}^{'}\right]_8$ denoting $\left[f^{'}{}_8(P^*)\right]_8\left[f_{8*}^{'}\right]_8$ denoting $[f_8]_8$ $(Q_8^*,$ $h_{*}^{'}$ denoting $h_*^{'}$ denoting $h_8^{'}$ $(Q_8^*)$, and $\left[g_*^{'}\right]_8$

Denoting $\left[\dot{g}\right]_8(P_8^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_8]_8 + [m_a]_8 = -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8 - \left[\dot{f}_{8*}\right]_8[p_8^*](a_8)M_8 - \left[\dot{f_{8*}}\right]_8 p_8^*(a_8)W_8$$

$$m_8(t,0) = [h_*]_8 \int_0^\infty \beta_8(a_8)m_8(t,a_8)da_8 + \frac{[h_*^8]_8}{[h^*]_8}[p^*]_8(0)\, W_{,8} \qquad (30)$$

$$\frac{dW_8}{dt} = [g_*^8]_8[Q_8^*]M_8$$

Integrating $[m_8]_8(t,a_8)$ with respect to a from 0 to ∞, we have

$$\frac{dM_8(t)}{dt} = \int_0^\infty [mt]_8(t,a_8)da_8$$

$$= -\int_0^\infty [m_a]_8\,(t,a_8)da_8 - \int_0^\infty \mu(a_8)m(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + \left[\dot{f}_{8*}\right]_8 \int_0^\infty p_8^*(a_8)da_8\right)M_8 - [f_{8*}^8]_8 \int_0^\infty p_8^*(a_8)da_8 W_8$$

$$= m_8(t,0) - \int_0^\infty \mu_8(a_8)\, m_8(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + \left[\dot{f_{8*}}\right]_8 P_8^*\right)M_8 - \left[\dot{f_{8*}}\right]_8 P_8^* W_{,8}$$

From (30), we arrive at the equations

$$\frac{dW_8(t)}{dt} = \int_0^\infty ([h_*]_8\beta_8(a_8) - \mu_8(a_8))m_8(t,a_8)da_8$$

$$-\left([f_{8*}] + [f_{8*}]_8 + [f_{8*}']_8 P_8^*\right) M_8 + \left(\frac{[h_*']_8}{[h_*]^8}\, P_8^*(0) - f_{8*}' P_8^*\right) W_8,$$

(31)

$$\frac{dW_8(t)}{dt} [g'_*]_8\, Q_8^*\, M_{.8}$$

By the mean value theorem,

$$\int_0^\infty ([h_{"*}]_8 \beta_8(a_8) - \mu_8(a_8)) m_8(t, a_8) da_8 = ([h_*]_8 \beta_8(\theta_8) - \mu_8(\theta_8)) M_{,8}$$

Where $0< \theta_8 = \theta_8(t) < \infty$.

And if a $^8_8$ [a $^8_8$]2<4 [a $_8$]2<4 a $_8$e,

Re $\lambda_8 [\lambda_8]_8 = \text{Re}\lambda_8 [\lambda_8]_8 = -\frac{8}{8} a_8 [\text{a2}\ ]_8 < 0$.

Let ⊿=max$_{8<u<\infty}$ |h $_*$ [h $_*$] $_8$β $_8$($a_8$) − $\mu_8(a_8)$|. Then|$n_8$(t)| ⊿∀t . Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium(ρ $^*$ (a), $\boldsymbol{Q}^*$) is asymptotically stable if either a $^8_8$ <4 a $_8$ e and [a $_8$] > 1⊿,or [a $^8_8$] [a $_8$]e and (a $_8$ − $\sqrt{(a^8_8)_8 - 4(a_8)_8 e}$ > 1⊿

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+8}\right)$ and β $_8$($a_8$) = ae-$^{ma}$ Then ⊿=[$\mu_0$] + 2 Simple calculation leads to

$$\int_0^{\infty} \beta_8(a_8)\mathrm{e}^{-\int_0^{\mathrm{a}} \mu_8(\tau_8)+d\tau_8}\mathrm{da} = \frac{1}{(\mathrm{m}+\mu_0+1)^8}+\frac{1}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_8 + [\mu_0]$ 0.5. Assume$[f_8](x_8)= K_8x_8$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_8(x_8)= [0x_8]/(2 + x_8)$; the functional responses $[[f_8](x) = [0]\, x_8/(1 + x_8)$and $h_8(x_8)$=$(1+\omega_8x_8)/(1 + x_8)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_8 + [\mu_0] + [\mathrm{K}][P_0])}{[m_8]_8 + (\mu_0)_8 + [\mathrm{K}][P_0])_8^3\ [\mathrm{K}]_8[P_0]_8^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_8$i e $[P_0] < b/(\delta_8 - b_8)$then the equilibrium$(P_0)$ $[P_0]_8(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_8$

$[f_{8*}] = K_8\ \mathrm{P}^*\ [f_{8*}] = \frac{\sigma \mathrm{Q}_8^*}{8+\mathrm{Q}_8^*}$ $[h_*] = \frac{8+\omega_8\mathrm{Q}_8^*}{8+\mathrm{Q}_8^*}$

$[f_{8*}] = K_8\ [\mathrm{f}_{8*}^{8}] = \frac{\sigma_8}{(8+\mathrm{Q}_8^*)^8}[\ \mathrm{g}_*^{8}] = \frac{\sigma_8}{(8+\mathrm{P}_8^*)^8}$ $[\mathrm{h}_*^{8}] = \frac{\omega_8-8}{(8+\mathrm{Q}_8^*)^8}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_8}{(m + [\mu_0]_8 + 1 + \frac{K_8b_8}{\delta_8 + b_8} + \sigma)^8} + \frac{1\omega_8}{([m_8]_8 + [\mu_0]_8 + 1 + \frac{K_8b_8}{\delta_8 + b_8} + \sigma_8)^3} < 1$$

And

$$\frac{1}{(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^8} + \frac{1\omega_8}{(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8} + (m_8 + [\mu_0]_8 + \frac{K_8 b_8}{\delta_8 - b_8} + \sigma_8)\, Q_8^*)^3$

$= m_8 + [\mu_0]_8 + 3 + \frac{K_8 b_8}{\delta_8 - b_8}(1 + Q^*)(1 + \omega_8 Q_8^*$

Now

$$\left[a_{11}\right] - 4\left[a_{88}\right] e = \left(\frac{1 K_8 b_8}{\delta_8 - b_8}\right)^8 + \frac{4 K_8 b_8 \sigma_8 Q_8^*}{(\delta_8 - b_8)(8 + Q_8^*)} + \frac{\sigma_8^8 Q_8^{*8}}{(1 + Q_8^*)^8}$$

$$\frac{4\, b_8 \sigma_8 (\delta_8 - b) Q^*}{\delta_8 (1 + Q_8^*)^8} - \frac{4\, (\delta_8 - b_8)^8 (1 - \omega_8) Q_8^*}{\delta_8 (1 + \omega_8 Q_8^*)(1 + Q_8^*)} P^*(0)$$

And

$$\Delta = \lim_{0 < a < \infty} |[h_*]_8\, a e_8^{mu}\ \ [\mu_0]_8 - \frac{a_8}{a_8 + 8}$$

Hence if $[a_8^8 > 4(a_8)_8 e$ and

$$[a_8] = \frac{1K_8 b_8}{\delta_8 - b_8} + \frac{\sigma_8 Q_8^*}{1 + Q_8^*} > 1\Delta$$

**The nontrivial equilibrium is stable.**

IF we define

$$[a_8]_8 = [f_{8*}]_8 + [f_{8*}]_8 + [f_{8*}] P_8^*$$

$$[a_8]_8 = \left[f_{8*}^{8}\right] P_8^* - \frac{h_*^8}{h_*} P^*(0)$$

$$e = [g_*^8]_8 Q_8^*$$

$$\eta_8(t) = \left[h_*\right]_8 \beta_8 \left(\theta_8(t)\right) - \mu_8(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\left[a_8\right] M_8 - [a_8]_8 W_8 - [n(t)] M_8$$

$$w_8(t) = eM_8$$

Write

$$A = \begin{pmatrix} -[a_8]_8 - [a_8]_8 \\ e \qquad 0 \end{pmatrix}$$

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| $B=(t)=\begin{pmatrix} [\eta(t)]_8 & 0 \\ 0 & 0 \end{pmatrix}$ |
| And |
| $X(t)=\begin{pmatrix} M_8 \\ W_8 \end{pmatrix}$ |
| The solution of (32)is |
| $X_8(t) = e^{A,t}[X]_8(0) + \int_0 e^{At}$ |
| This equation notwithstanding its importance in all the models stand out for its importance for the following models<br><br>to which it is applied: |
| **Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016** |

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| (217)These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field. |
| (218)The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets. |
| (219) Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state. |
| (220) Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands. |
| (221)The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI |
| (222) The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9]. |
| (223)This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive. |

(224)More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

(225)Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.

(226)Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.

(227) The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.

(228)This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.

(229)The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.

(230)This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.

(231) For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.

(232)The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(233)Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.

(234)We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.

(235)When this is true—as is often the case—the system serves as (=) a "bath" for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called **thermalization.**

(236)However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.

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| (237)Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1). |
| (238)Bringing them to thermal equilibrium requires (e) an external bath. |
| (239) But so far, there has been no rigorous proof of (e) the existence of **such many-body localization**. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption. |
| (240)The concept of **localization** was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations). |
| (241)The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century. |
| (242)But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information. |
| (243)Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment. |
| (244)For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood. |
| (245) It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7]. |
| (246) Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature. |
| (247) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators. |
| (248)However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8]. |
| (249)Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems. |
| (250) Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL. |
| (251)As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked. |
| (252)His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice. |
| (253) An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site. |

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| (254)This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself). |
| (255)This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition. |
| (256)However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime: |
| (257) These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects. |
| (258)It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies). |
| (259) This is a realistic scenario for (e) all known systems. |
| (260)On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects. |
| (261)On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other. |
| (262)No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true. |
| (263) His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found. |
| (264)Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization? |
| (265)Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects? |
| (266)Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions. |
| (267)So in those systems, the quantum phase transition does indeed exist. |
| (268)But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension. |
| (269)Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters. |

# PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION ATTRITION MODELS

## MODULE NINE

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Model is applied to following formulations Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society**

**Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP This paper is divided into four parts. In part A, some general considerations about (e&eb) gravitational radiation are followed by (e) a treatment of the scalar wave equation in the manner later to be applied to (e&eb) Einstein's field equations. In part B, a co-ordinate system is specified which is suitable for (e) investigation of outgoing gravitational waves from (e) an isolated axi-symmetric reflexion-symmetric system. The metric is expanded in (eb) negative powers of (e) a suitably defined radial co-ordinate r, and the vacuum field equations are investigated in detail. It is shown that the flow of information to infinity is controlled by (e&eb) a single function of (e) two variables called the news function. Together with initial conditions specified on a light cone, this function fully defines (eb) the behaviour of the system. No constraints of (e) any kind are encountered. In part C, the transformations leaving (eb) the metric in the chosen form are determined. An investigation of the corresponding transformations in (eb) Minkowski space suggests (eb) that no generality is lost by (e) assuming that the transformations, like the metric, may be expanded in (eb) negative powers of r. Gravitational Waves in General Relativity VII Waves from Axi-Symmetric Isolated Systems Royal Society Proceedings H. Bondi, M. G. J. van der Burg, A. W. K. Metzner Published 21 August 1962.DOI: 10.1098/rspa.1962.0161 STOP There is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant Λ, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive Λ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the Λ>0 generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible. Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of (e) GW signals from BBHs with component spins $\chi_{1,2} \in [-0.99, 0.99]$ aligned (e&eb) with the orbital angular momentum, component masses $m_{1,2} \in [2,48] M_\odot$, and (e&eb) total mass $M_{total} \leq 50\ M_\odot$. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin**

**bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins >+0.6 and component masses ≳20 M⊙. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201**

## 27. MODEL FORMULATION

Let $p_9(t, a_9)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(g_9)Q_9$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_9(0) = 0, \quad g_9(P_9) > 0, \ P_9 \ 0, \qquad \lim_{P_9 \to \infty} g_9(P_9) = (K_9)_9 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_9 + (\rho_a)_9 + D_9(a_9, P_9, Q_9)\rho_9 + 0, \qquad 0 < a_9 \ \infty,$$

$$\rho_9(t,0) = \int_0^\infty n_9(a_9, P_9, Q_9)\rho_9(t_9, a_9)da_9$$

Where $D_9$ is the mortality function and $n_9$ is the fertility function.

Put $_9$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_9}{dt}\right]_9 = x_9\left[\frac{dQ_9}{dt}\right]_9 - x_9\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_9}{dt}\right]_9 = x_3\left[\frac{dQ_3}{dt}\right]_9 - x_4\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_9}{dt}\right]_9 - x_6\left[\frac{dQ_3}{dt}\right]_9$$

Where $x_9, x_9, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_9$ and $g_9(P_9)Q_9$

Now when you add the term $(b_9Q_9)$ with $g_9(P_9)(Q_9)$ we are infect see is that $(b_9Q_9)$ is credited incrementally part of $g_9(P_9)(Q_9)$ or instant nearly, when subtracted from $\frac{dQ_9}{dt}$ , it means that $g_9(P_9)(Q_9)$ is slowly eaten by $\frac{dQ_9}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_9 = \mu_9(a_9) + (f_9)_9(P_9)_9 + (f_9)_9(Q)_9 \quad (2)$$

$$n_9 = \beta_9(a_9)h_9(Q_9), \quad (3)$$

Where $\mu_9$the age- dependent debit rate whit is $\mu_9(a_9)$ 0 $\forall a > 0$; $(f_9)_9$ is the density- dependent death rate with $f_9(0) = 0$, $(f_9^9)_9(P_9)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_9)_9(P_9) = (k_9)_9$ ; $f_9$ is the functional response of creditors to prey with $[f_9]_9(0) = 0, (f_9^9)_9(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_9)_9(Q_9) = [k_9]$ ; $\beta$ is the age dependent birth

rate with $\beta(a) \geq 0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h'(Q) \leq 0, \forall Q \geq 0$, and $\lim_{Q\to\infty} h(Q) = k_3 \geq 0$.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_9 + (\rho_a)_9 = -\mu_9(a_9) + (f_9)_9 P_9) + \big(f_9)_9(Q_9)\big)(\rho_9)_9 \tag{4}$$

$$\rho_9(t_9, 0) = h_9(Q_9) \int_0^\infty \beta_9(a_9)\rho_9(t, a_9)da, \tag{5}$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(P_9)Q_9 \tag{6}$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_9$ and $Q_9$ of (4-6) satisfy

$$0 \leq \rho_9(t, a_9) \leq \rho_9^*(t, a_9),$$

$$0 \leq Q_9(t) \leq Q_9^*(t),$$

Where $\rho_9^*(t, a_9)$ and $Q_9^*(t)$, are solutions of

$$(\rho_t^*)_9 + (\rho_a^*)_9 = -(\mu_9(a_9) + (f_9)_9(P_9^*) + (f_9)_9 Q_9^*))\, \rho_9^*, \tag{7}$$

$$\rho_9^*(t, 0) = \int_0^\infty \beta_9(a_9)\, \rho_9^*(t, a_9)da_9 \tag{8}$$

$$\frac{dQ_9^*}{dt} = -b_9 Q_9^* + g_9(P_9^*)Q_9^*, \tag{9}$$

With $P_9^*(t) = \int_0^\infty \rho_9^*\ (t, a_9)da_9$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_9^*(t)}{dt} = P_9^*(t)\big(r_9 - (f_9)(P_9^*) - (f_9)_9(Q^*)\big), \quad (10)$$

$$\frac{dQ_9^*(t)}{dt} = Q_9^*(t)(-b_9 + g_9(P_9^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_9(a)e^{-\int_0^a \mu_9(\tau)dt}e^{-ra}\ da_9 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_9^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(t)dt}\ da_9 > 1, \quad (H1)$$

Which guarantees $r_9 > 0$.

*Remark.* If $\beta_9(a) = [\beta_0]_9 a_9 e^{-\alpha_9 a_9}, \mu_9(a_9) \equiv \mu_9$ then the solution of (12) is

$$r_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_9]_9(P_9) = 0, [f_9](Q_9) = \lambda_9 Q_9\ h_9(Q_9) = 1, \quad and\ \ g_9(P_9) = c_9 P_9$$

$$[\rho_9]_9 + [\rho_a]_9 = -(\mu_9 + \lambda_9 Q_9)\rho_9$$

$$\rho_9(t,0) = \int_0^{\infty} [\beta_0]_9 a_9 e^{-\alpha_9 a_9} \rho_9(t,a_9)\, da_9$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + c_9 P_9 Q_9$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_9}{dt} = \omega_9 P_9 - \lambda_9 P_9 Q_9$$

$$\frac{dQ_9}{dt} = b_9 Q_9 - c_9 P_9 Q_9$$

Where $\omega_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$

THE EQUILIBRIUM $(\rho_0)_9(a_9), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_9(a_9), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_9 + (\rho_a)_9 = -\big(\mu_9(a_9) + (f_9)_9(P_9)\big)\rho,$$

$$\rho_9(t,0) = \int_0^{\infty} \beta_9\,(a_9)\rho_9(t,a_9)da_9 \qquad (13)$$

The solution of (13) has the form of

$$\rho_9(t,a_9) = Ce_9^{r(t-a_9)} - \int_0^a (\tau)d\tau - \int_0^9 (f_9)_9\big(P_9(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_9(t) = C_9 e^{r_9 t - \int_0^t (f_9)_9(P_9(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_9^{-9}(r)$ if $k_9 > r$. This is equivalent to

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{(k_9)_9 a_9} da_9 < 1.$$

In fact, if we define

$$R_9(s_9) = \int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau)d\tau} e^{-s_9 a_9}\, da_9$$

$R(s)$ is a decreasing function. $R_9(r_9) = 1$ And $R_9(k_9)_9 < 1$ imply $r_9 < (k_9)_9$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_9(a_9), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{k_9 a_9} da_9 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_9 \equiv \int_0^\infty (\rho_0)_{99}(a_9)da_9$ *satisfies*

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{-(f_9)_9(P_0)a_9} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_9(a_9), 0)$

Let $u_9(t,a_9) = \rho(t,a_9) - \rho_0(a_9)$ Then

$$N_9(t) \equiv \int_0^\infty u_9\,(t, a_9)da_9 = \int_0^\infty \rho_9(t, a_9)da_9 - \int_0^\infty (\rho_0)_9(a_9)da = P_9(t) - P_0)_9$$

And

$$(u_9)_9 + (u_a)_9 = (\rho_t)_9 + (\rho_a)_9 - (\rho_0^9)_9(a_9)$$

$$= -\mu_9(a_9) + (f_9)_9(P_9) + (f_9)_9(Q_9))\,(u_9 + (\rho_0)_9) + \mu_9(a_9)(\rho_0)_9 + (f_9)_9(P_0)_9(\rho_0)_9$$

$$= -\mu_9(a_9)u_9 - (f_9)_9(P_9) + (f_9)_9(Q_9))\,u_9 - (\rho_0)_9\,(f_9)_9(P_9) - (f_9)_9(P_0)_9$$

$$+ (f_9)_9(Q_9)), \qquad (15)$$

Where $(\rho_0^9)_9(a_9)$ denotes $d(\rho_0)_9(a_9)/da_9$

Expansion $(f_9)_9(P_9)$ and $(f_9)_9(Q_9)$ at $(P_0)_9$ and 0 yield respectively

$$(f_9)_9(P_9) = (f_9)_9((P_0)_9 + N_9)$$

$$= (f_9)_9(P_0)_9 + (f_9^9)_9(P_0)_9N_9 + o_9(N_9), \qquad (16)$$

$$(f_9)_9(Q_9)$$

$$= (f_9^9)_9(0)Q + o_9(Q_9) \qquad (17)$$

Write $f_9(P_0) = f_{90}, f_9'(P_0) = f_{90}'$, and $f_9'(0) = f_{90}'$ and substitute (16) and (17) into (15) to obtain

$$(u_9)_9 + (u_a)_9 \approx -(\mu_9(a_9) + (f_{90})_9)u - (f_{90}^9)_9(\rho_0)_9(N_9) - (f_{90}^9)_9(\rho_0)_9(Q_9)$$

$$u_9(t, 0) = \rho_9(t, 0)(\rho_0)_9(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t, a) + (\rho_0)_9(a_9)da_9 - (\rho_0)_9\,(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t, a_9)da_9 + (h_9(Q_9) - 1)(\rho_0)_9\,(0) \qquad (18)$$

Expanding

$$h_9(Q_9) = 1 + (h^9)_9(0)Q_9 + 0(Q_9) \equiv 1 + h_0'Q_9 + 0Q_9$$

And substituting it into (18), we arrive at

$$u_9(t,0) \approx \int_0^\infty \beta_9(a_9)u_9(t,a_9)da_9 + (h'_{a_9})(\rho_0)_9Q_9$$



PREDATOR-PREY POPULATION MODELS

The linearization is the following :

$$[u_9] + [u_9] = -(\mu_9(a_9) + [t_{90}]_9u_9 - [\alpha_9]_9(a_9)N_9 - [\alpha_9]_9(a_9)Q_9$$

$$u_9(t,0) = \int_0^\infty \beta_9(a_9)u_9(t,a_9)da_9 + [\xi_9]Q_9 \qquad (19)$$

$$\frac{dQ_9}{dt} = -b_9Q_9$$

where

$$[\alpha_9]_9(a_9) = [t_9^9a]_9[p,o]_9(a_9)d2(a_9) = [f_{90}^9]_9(\rho\, o\,)_9(a_9)\xi = [\mathrm{h\ o}^9]_9[\rho\mathrm{o}]_9(\mathrm{o}) \text{ and } \mathrm{b}_9 = \mathrm{b}_9 - g_9(P_9)$$

If $g_9(P_0) > b$, clearly, $(\rho_0)_9(a_9), 0)$ is unstable and the predator population is persistent. Assume $g_9(P_0)_9 < b$.

Then, $Q_9(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_9 + u_a = -(\mu_9(a_9) + (f_{90})_9)u_9 - (\alpha_9)_9(a_9)N_9$$

$$= -(\mu_9(a_9) + f_{90})u_9 - \alpha_9(a_9)\int_0^\infty u_9(t,a_9)da_9$$

$$u_9(t,0) = \int_0^\infty \beta_9\,(a_9)u_9(t,a_9)da_9$$

Using the method developed by Cushing [2, 3], we suppose

$$u_9(t,a_9) = \tilde{u}_9(a_9)^{e^{\gamma(t-a_9)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_9}\,da_9$. Then,

A solution of the first differential equation for $\tilde{u}(a_9)$ gives,

$$\tilde{u}_9(a_9) = e^{-\int_0^a \mu_9(\tau_9)d\tau-(f_{90})_9 a}\left(\tilde{u}_9(0) - w\int_0^a e^{\int_0^\tau \mu_9(s_9)ds+(f_{90})}\tau_9\,(\alpha_9)_9(\tau_9)e^{\gamma_9\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\right)\tilde{u}_9(0)$$

$$+w\int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\int_0^a (\alpha_9)_9(\tau_9)e^{-\gamma_9\tau_9}e^{(f_{90})_9\tau_9+\int_0^\tau \mu_9(s_9)ds}d\tau_9\,da_9$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}da\right)(\tilde{u}_9(0))_9$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\int_0^a e^{(f_{90})_9\tau_9+\int_0^\tau \mu_9(s_9)ds+y\tau}\alpha_9(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(t_9)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_9\,(a_9)Y_9(\gamma_9,a_9)da_9\right)u_9(0)$$

$$+w\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)\int_0^a(\alpha_9)_9(\tau)_9(Y^{-9})_9(\gamma_9,\tau_9)d\tau_9\,da_9=0,$$

(24)

$$\left(\int_0^{\infty}Y_9(\gamma_9,a_9)da_9\right)\tilde{u}_9(0)$$

$$+\left(-1-\int_0^{\infty}Y_9(\gamma_9,a_9)\int_0^a(\alpha_9)_9(\tau)_9(Y^{-9})_9(\gamma_9,\tau_9)d\tau_9\,da_9\right)w=0.$$

Equation (20) has a solution of the form of $u_9(t,a_9)=u_9(a_9)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_9(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_9(\gamma_9,a_9)=\int_0^a(\rho_0)_9(\tau_9)Y^{-9}\,(\gamma_9,a_9)d\tau_9$. We obtain the characteristic equation

$$f_{90_9}^{l}\left(\int_0^{\infty}Y_9(\gamma_9,a_9)da_9\right)\left(\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)I_9(\gamma_9,a_9)da_9\right)$$

$$+\left(1-\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)da\right)\left(1+(f_{90}^{l})_9\int_0^{\infty}Y_9(\gamma_9,a_9)I_9(\gamma_9,a_9)da_9\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_9(a_9) = (\rho_0)_9(0)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - (f_{90})_9 a_9}$$

$$I_9(\gamma_9, a_9) = \int_0^a (\rho_0)_9(\tau_9)e^{\gamma_9\tau_9 + (f_{90})_9\tau + \int_0^\tau \mu_9(s_9)ds_9}d\tau$$

$$= \frac{(\rho_0)_9(0)}{\gamma_9}(e^{\gamma_9 a_9} - 1)$$

Let

$$E_9(a_9) = e^{-(f_{90})_9 a_9 - \int_0^a \mu_9(\tau_9)d\tau_9}$$

Then $[Y_9]_9[\gamma_9]_9, a_9) = E_9(a_9)e^{-\gamma_9 a_9}$. by a simple calculation,

$$\int_0^\infty \beta_9(a_9)Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma}\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}\ da\right);$$

$$\int_0^\infty Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma_9}\left(\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}\ da_9\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_9(0)}{\gamma_9}[f_{90}^9]\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}da_9\ (\int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9)$$

$$+\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right).$$

$$\left(1+\frac{(\rho_0)_9(0)}{\gamma_9}(f_{90}^{'})_9\left(\int_0^{\infty}E_9(a_9)da_9-\int_0^{\infty}E_9(a_9)e^{-\gamma_9 a_9}da_9\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right)(1+\frac{(f_{90}^{'})_9}{\gamma_9}(\rho_0)_9(0)\int_0^{\infty}E_9(a_9)da_9)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (94), to

$$1-\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)da_9=1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\neq 0,$$

Since $\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)\int_0^{a}\alpha_9(\tau_9)Y^{-9}(\gamma_9,\tau_9)d\tau_9 da_9>0$. Therefore, the characteristic solution is

$$\gamma_9=-(f_{90}^{'})_9(\rho_0)_9\int_0^{\infty}E_9(a_9)da_9.$$

Since $f_9^{'}(P_0)_9>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_9),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_9),0)$ which is related to the extinction of predators (debit summation types: that which acts as a resource)and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $( p^*(a), Q^* )$. It must satisfy the equations:

$$[p_a^*]_9 + (\mu_9(a_9) + [f_9]_9( P_9^*) + [f_9]_9(Q_9^*) )P_9^* = 0, \quad (26)$$

$$p_9^*(0) = h_9(Q_9^*) \int_0^\infty \beta_9(a_9)p_9^*(a_J)p_9^*(a_9)da_9 \quad (27)$$

$$g_9(P_9^*) - b_9 = 0, \quad (28)$$

Where $P_9^*=\int_0^\infty p_9^*(a_9)\, da_9$

Equation (28) yields $P^* = g^{-9}(b)$,

From (26),

$$p_9^*(a) = p_9^*(0)e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}$$

Substituting into (27), we need to solve

$$\left(1 - h_9(Q_9^*)\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}da_9\right)p_9^*(0) = 0$$

For $Q_9^*$

Fix $p_9^*(0)$ and $P_9^*$ Let

$$H_9(z_9) = 1 - h_9(z_9)\int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9}\, da_{,9}$$

And assume that

$$[k_3]_9 \int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(g_9^{-9}(b_9))a_9 - [k_9]_9 a_9}\, da_9 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_9(z_9)] > 0.$$

Clearly,

$$\frac{dH_9}{dz_9} = h_9(z)\left([f_9']_9(z_9)\right)\int_0^\infty a_9\beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - f_9(P_9^*)a_9 - f_9(z_9)a_9} da_{\square}$$

$-h'(z_9)\int_0^\infty \beta_9(a_9)\, e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9}\ da_9$

y the assumptions on $\left[f_9^9\right]_9 and\ [h^9]^9 dH_9/dz > 0$. Hence, if

$$H_9(0) = 1 - \int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9} da_9 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_9(0) < 0$ follows if $g_9^-(b_9) < [f_9^{-9}]_9(r_9)$, where $r_9$ is the solution of (12).

Indeed, since

$$R_9(s_9) = \int_0^{\infty} \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 e^{-y\frac{a}{9}} da_9}$$

Is a decreasing function of $s_9$ when $P_9^* = [g^{-9}]_9(b_9) < [f_9^{-9}]_9(r_9)$, that is $[f_9]_9(P_9^*) < r$, $R\big(f_9(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_9(g^{-9}(b))) > 1, i.e.,$

$\int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9\left(g_9^{-9}(b)\right)a_9} > 1.$ (H4)

Then the system (4)-(6) has a unique positive equilibrium $(P_9^*(a_9), Q_9^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_9^*(a_9), Q_9^*)$ is more complicated. We again consider perturbations $m_9(t_9 a_9) = p_9(t_9, a_9) - p_9^*(a_9)$ and $W_9(t) = Q_9(t) - Q^*$.Then

$M_9(t) \equiv \int_0^{\infty} m_9(t, a_9) da_9 = P_9(t) - P_9^*,$

$$[m_9]_9 + [m_a]_9 = [p_9]_9 + [p_a]_9 - \left[p^*\right]_9(a_9) \approx -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9$$

$$-[f_{9*}^9]_9 p_9^*(a_9)M - [f_{9*}^9]_9 p_9^*(\mathrm{a})W_9,$$

$$m_9(t,0)p_9(t,0) - p_9^*(0) \approx h_* \int_0^{\infty} \beta_9(a_9) m_9(t, a_9)\, da + \frac{[h_*^9]_9}{[h_*]_9}\, p_9^*(0)W_9$$

$$\frac{dW_9}{dt} = \frac{dQ_9}{dt} \approx [g_*^9]_9 Q_9^* M_{,9}$$

Where $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $\left[f_{9*}^{'}\right]_9$ denoting $\left[f^{'}{}_9(P^*)\right]_9 \left[f_{9*}^{'}\right]_9$ denoting $[f_9]_9$ $(Q_9^*,\quad h_9^*$ denoting $h_*^{'}$ denoting $h_9^{'}$ $(Q_9^*)$, and $\left[g_*^{'}\right]_9$

Denoting $\left[g^{'}\right]_9(P_9^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_9]_9 + [m_a]_9 = -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9 - \left[f^{'}{}_{9*}\right]_9 [p_9^*](a_9)M_9 - \left[f_{9*}^{'}\right]_9 p_9^*(a_9)W_9$$

$$m_9(t,0) = [h_*]_9 \int_0^\infty \beta_9(a_9)m_9(t,a_9)da_9 + \frac{[h_*^9]_9}{[h^*]_9}[p^*]_9(0)\,W_{,9} \qquad (30)$$

$$\frac{dW_9}{dt} = [g_*^9]_9[Q_9^*]M_9$$

Integrating $[m_9]_9(t,a_9)$ with respect to a from 0 to ∞, we have

$$\frac{dM_9(t)}{dt} = \int_0^\infty [mt]_9(t,a_9)da_9$$

$$= -\int_0^\infty [m_a]_9\,(t,a_9)da_9 - \int_0^\infty \mu(a_9)m(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + \left[f^{'}{}_{9*}\right]_9 \int_0^\infty p_9^*(a_9)da_9\right)M_9 - [f_{9*}^9]_9 \int_0^\infty p_9^*(a_9)da_9 W_9$$

$$= m_9(t,0) - \int_0^\infty \mu_9(a_9)\, m_9(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + \left[f_{9*}^{'}\right]_9 P_9^*\right)M_9 - \left[f_{9*}^{'}\right]_9 P_9^* W_{.9}$$

From (30), we arrive at the equations

$$\frac{dW_9(t)}{dt} = \int_0^{\infty} ([h_*]_9\beta_9(a_9) - \mu_9(a_9))m_9(t,a_9)da_9$$

$$-\left([f_{9*}] + [f_{9*}]_9 + \left[f_{9*}^{'}\right]_9 P_9^*\right) M_9 + \left(\frac{\left[h_*^{'}\right]_9}{[h_*]^9}\ P_9^*(0) - f_{9*}^{'} P_9^*\right) W_9,$$

(31)

$$\frac{dW_9(t)}{dt}\left[g^{'}{}_*\right]_9 Q_9^*\, M_{.9}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{"*}]_9\beta_9(a_9) - \mu_9(a_9))m_9(t,a_9)da_9 = ([h_*]_9\beta_9(\theta_9) - \mu_9(\theta_9))M_{,9}$$

Where $0< \ \theta_9 = \theta_9(t) < \infty$.

And if $a_9^9[a_9^9]2<4\ [a_9]2<4\ a_9e$,

$\text{Re}\,\lambda_9[\lambda_9]_9 = \text{Re}\lambda_9[\lambda_9]_9 = -\frac{9}{9}a_9[a2\ ]_9 < 0.$

Let $\Delta=\max_{9<u<\infty}\left|h_*[h_*]_9\beta_9(a_9) - \mu_9(a_9)\right|$. Then$|n_9(t)|$ ⍰⊿∀t ⍰0. Hence we obtain the following result.

**THEOREM 4**

Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_9^9 <4\ a_9e$ and $[a_9] > 1\Delta$,or $[a_9^9]$ ⍰4$[a_9]e$ and $(a_9 - \sqrt{(a_9^9)_9 - 4(a_9)_9e) > 1\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+9}\right)$ and $\beta_9(a_9) = ae^{-ma}$ Then $\Delta=[\mu_0] + 2$ Simple calculation leads to

$$\int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9) + d\tau_9 da} = \frac{1}{(m + \mu_0 + 1)^9} + \frac{1}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_9 + [\mu_0]$ 0.5. Assume$[f_9](x_9)= K_9 x_9$ Then the hypothesis (H9) is satisfied We take the conversion rate function $g_9(x_9)= [0x_9]/(2 + x_9)$; the functional responses $[[f_9](x) = [0]\, x_9/(1 + x_9)$and $h_9(x_9)=(1+\omega_9 x_9)/(1 + x_9)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_9 + [\mu_0] + [K][P_0])}{[m_9]_9 + (\mu_0)_9 + [K][P_0])_9^3\ [K]_9[P_0]_9^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_9$i e $[P_0] < b/(\delta_9 - b_9)$then the equilibrium$(P_0)$ $[P_0]_9(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_9$

$[f_{9*}] = K_9\ P^*\ [f_{9*}] = \frac{\sigma Q_9^*}{9+Q_9^*}$ $[h_*] = \frac{9+\omega_9 Q_9^*}{9+Q_9^*}$

$[f_{9*}] = K_9\ [f_{9*}^9] = \frac{\sigma_9}{(9+Q_9^*)^9}\ [g_*^9] = \frac{\sigma_9}{(9+P_9^*)^9}$ $[h_*^9] = \frac{\omega_9 - 9}{(9+Q_9^*)^9}$

In order to have (H3) and (H4) hold,$m, \mu_0, K, b, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_9}{(m+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9+b_9}+\sigma)^9}+\frac{1\omega_9}{([m_9]_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9+b_9}+\sigma_9)^3}<1$$

And

$$\frac{1}{(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9})^9}+\frac{1\omega_9}{(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9}+(m_9+[\mu_0]_9+\frac{K_9 b_9}{\delta_9-b_9}+\sigma_9)\,Q_9^*)^3$$

$$=m_9+[\mu_0]_9+3+\frac{K_9 b_9}{\delta_9-b_9}(1+Q^*)(1+\omega_9 Q_9^*$$

Now

$$\begin{bmatrix}1\\1\end{bmatrix}-4\begin{bmatrix}a\\9\end{bmatrix}e=\left(\frac{1K_9 b_9}{\delta_9-b_9}\right)^9+\frac{4K_9 b_9\sigma_9 Q_9^*}{(\delta_9-b_9)(9+Q_9^*)}+\frac{\sigma_9^9 Q_9^{*9}}{(1+Q_9^*)^9}$$

$$\frac{4\,b_9\sigma_9(\delta_9-b)Q^*}{\delta_9(1+Q_9^*)^9}-\frac{4\,(\delta_9-b_9)^9(1-\omega_9)Q_9^*}{\delta_9(1+\omega_9 Q_9^*)(1+Q_9^*)}P^*(0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_9 \, ae^{mu}_9 \quad [\mu_0]_9 - \frac{a_9}{a_9+9}$$

Hence if $[a^9_9 > 4(a_9)_9 e$ and

$$[a_9] = \frac{1K_9 b_9}{\delta_9 - b_9} + \frac{\sigma_9 Q^*_9}{1 + Q^*_9} > 1\Delta$$

Some of the examples to which the model has been applied are:

**New look at the galaxy-shaping power of black holes**

(239) "It creates (eb) these bubbles, and the bubbles keep the plasma hot.

(240) That's what prevents galaxies from (e) becoming even bigger than they are now."

(241) Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.

(242) In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.

(243) Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.

(244) "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.

(245) "It's like nature's thermostat, which keeps these galaxies from (e) growing.

(246) If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."

**A new look at the galaxy-shaping power of black holes**

(247) An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the

(248) McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.

(249) Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.

**Doomed satellite missions**

(250) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.

(251) At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.

(252) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.

(253) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.

(254) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.

(255) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people

(256) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."

(257) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.

(258) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(259) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.

(260) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.

(261) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.

(262) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.

| |
|---|
| (263) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well. |
| (264) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma. |
| (265) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6. |
| (266) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus. |
| (267) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (268) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (269) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (270) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (271) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (272) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (273) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (274) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (275) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (276) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (277) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (278) We know dark matter exists only from (e) its gravitational effects. |
| (279) It holds together galaxies and (e&eb) galaxy clusters. |
| (280) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (281) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |

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| (282) | They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (283) | X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (284) | They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (285) | This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (286) | An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |
| (287) | Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (288) | "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (289) | The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (290) | Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (291) | Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (292) | "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (293) | Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (294) | "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (295) | Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) |

consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories.

(296) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen.

(297) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer.

(298) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants.

(299) These were treated as (=) "aware" and "unaware" trials, respectively.

(300) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it.

(301) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another.

(302) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion.

(303) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness.

(304) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said.

(305) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness."

(306) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner.

(307) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain.

(308) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said.

(309) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

(310) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter.

(311) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined.

(312) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions.

(313) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.

(314) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.

(315) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.

(316) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.

(317) Noda and his colleagues instead studied the change in emissions over (e&eb) time.

(318) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.

(319) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).

(320) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.

(321) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.

(322) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(323) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking "palm tree" organization.

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| (324) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society |
| (325) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible? |
| (326) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem. |
| (327) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other. |
| (328) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible. |
| (329) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges. |
| (330) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century. |
| (331) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching. |
| (332) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom. |
| (333) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth. |

**Spheres on trees**

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| (334) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration. |
| (335) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement. |
| (336) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum. |
| (337) By visualizing the problem in this way, the researchers noticed that these particular graphs don’t (e) have lots of branches extending from lots of (e) other branches. |
| (338) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure. |

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| (339) The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network. |
| (340) As a result, it provides (eb) an important clue for finding the global minimum. |
| (341) It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum. |
| (342) This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large. |

**Thompson's problem 2**

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| (343) By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society |
| (344) The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima. |
| **(345)** The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node. |

**Implications for Thomson's and Smale's problems**

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| (346) Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before. |
| (347) "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around." |
| (348) As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem. |
| (349) "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view." |
| (350) In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges. |
| (351) From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear. |
| (352) "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK. |
| (353) "For larger numbers, there will likely be (=) multiple funnels. |
| (354) Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum." |

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| (355) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem. |
| (356) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said. |
| (357) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org |

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| The nontrivial equilibrium is stable. |
| IF we define $[a_9]_9 = [f_{9*}]_9 + [f_{9*}]_9 + [f_{9*}]\,P^*_9$ |
| $$[a_9]_9 = \left[f^{9}_{9*}\right]P^*_9 - \frac{h^9_*}{h_*}P^*(0)$$ |
| $$e=[g^9_*]_9\,Q^*_9$$ |
| $$\eta_9(t) = \left[h_*\right]_9\beta_9\left(\theta_9(t)\right) - \mu_9\left(\theta(t)\right),$$ |
| The system(32)can be written as the following system: |
| $$[M(t)]\left[a_9\right]M_9 - [a_9]_9W_9 - [n(t)]M_9$$ |
| $$w_9(t) = eM_9$$ |
| Write |
| $$A=\begin{pmatrix}-[a_9]_9 - [a_9]_9 \\ e \qquad 0\end{pmatrix}$$ |

$$B=(t)=\begin{pmatrix} [\eta(t)]_9 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_9 \\ W_9 \end{pmatrix}$$

The solution of (32)is

$$X_9(t) = e^{A,t}[X]_9(0) + \int_0 e^{At}$$

**Note: There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and**

**sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by**

**term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation**

**alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation,**

**appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module**

**LEN NUNNEYS**

## Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

## Module numbered one

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems

and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1965); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be referred to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 6985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_6}{dt} = [N(t)]_6[(F_6)(R(t)]_6 - [M]_6[R(t)]_6$$

$$\frac{d[R(t)]_6}{dt} = [B]_6\,[(R)_6(t-T_6)] - [D]_6[(R(t)]_6 - [N(t)]_6\left[G\left(R(t)\right)_6\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_6}{dt} = a_6[N(t)]_6 - a_6[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_6 = a_3[N(t)]_3 - a_4[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_6 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

irst, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.6b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The

regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 6a is at the

**TIME DELAY**

FIG. 6. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.6$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**DELAY DUE TO PREDATOR MATURATION**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_6}{dt} = N_6(t - T_6)F_6(R_6(t - T_6) - N_6(t)M_6((R_6(t) \qquad (5)$$

$$\frac{dR_6(t)}{dt} = B_6\big(R_6(t) - D_6(R_6(t)\big) - N_6(t)G_6(R_6(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_6}{dR_6} + \frac{N_6 dG_6}{dR_6} - \frac{dB_6}{dR_6} > \frac{N_6 G_6}{2M_6}\left[\frac{dF_6}{dR_6} + \frac{dM_6}{dR_6}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_6/dR_6$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_6$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_6$) and the other derived from the non-delayed or "immediate" processes($I_6$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

ABSOLUTE STABILITY

$$L_6 = \frac{N_6 G_6}{2M_6}\frac{d(6nF_6)}{dR_6} (\geq 0) \qquad (6a)$$

$$I_6 = -\frac{N_6 G_6}{2M_6}\frac{d(6nM_6}{dR_6} (\geq 0) \qquad (6b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_6 > L_6 - I_6$, must be modified ahe $dM_6/dR_6$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_6 > L_6 - I_6, \quad \text{when} \quad P_6^2/4 > I \qquad (8a)$$

$$P_6 > L_6 - I_6 + \frac{(I_6 - P_6^2}{(L_6 - I)} \quad \text{when} \quad P_6^2/4 < I. \quad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

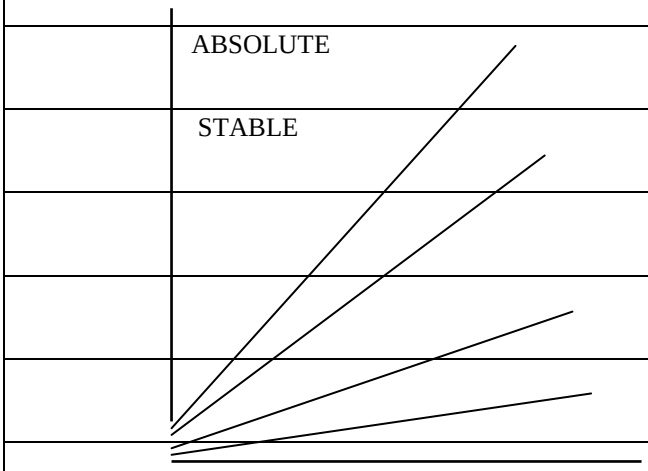



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

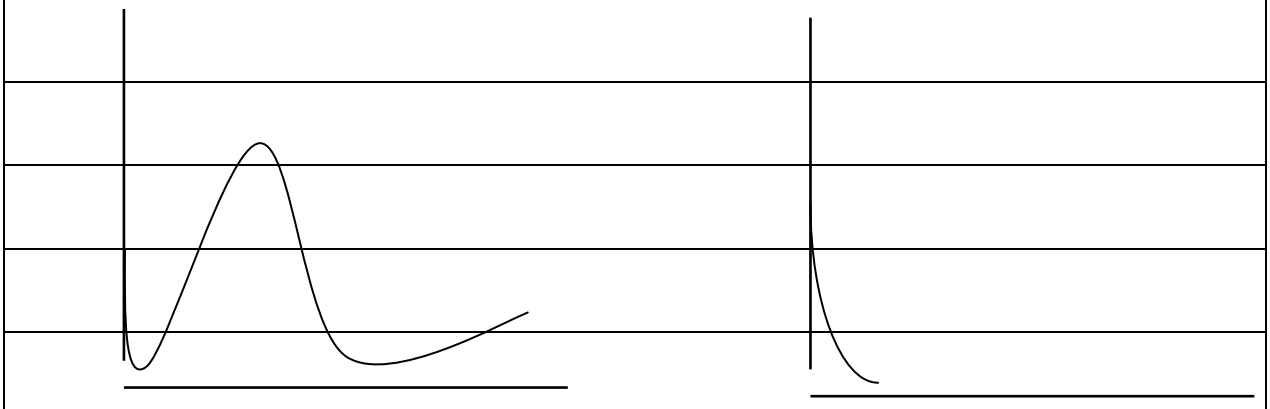
Can never completely offset the lagged effect,$L_6$; condition (8b) ensures that the total debit regulation, $P_6$, must always be positive even when $I_6$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_6$ and $L_6$. As $I_6$ increases relative to $L_6$ the zone of stability is enlarged; however,even when $I_6$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_6$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



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Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (6). Parameters in each case M=6; $dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6 = 0.05$ For (a) $NGdF/dR = 0.3, -N_6G_6dM_6/dR_6 = 0.6$; for (b)$NG_6dM_6/DR_6 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_6t}{dt} = N_6(t - T_6)F_6\big(R_6(t - T_6)\big) - N_6(t)M_6\big(R_6(t)\big)$$

(9)

$$\frac{dR_6t}{dt} = B_6\big(R_6(t - T_6)\big) - D_6\big(R_6(t)\big) - N_6(t)G_6\big(R_6(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) and the effective prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we

could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

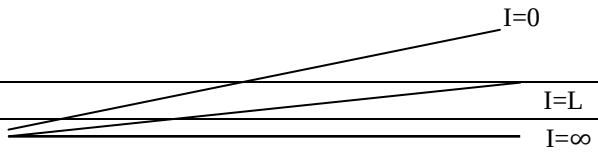
Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_6/dR_6$ is a measure of the control acting through the debit delay and setting $dB_6/dR_6 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_6/dR_6$ is a measure of the control acting through the credit delay although setting $dF_6/dR_6 = 0$ does not negate the effect of a predator delay; the value of $M_6$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_6/dR_6$and normalized by $M_6$ (see6a). The absolute stability of such a system is summarized in Fig. 6 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

TOTAL

ABSOLUTELY STABLE

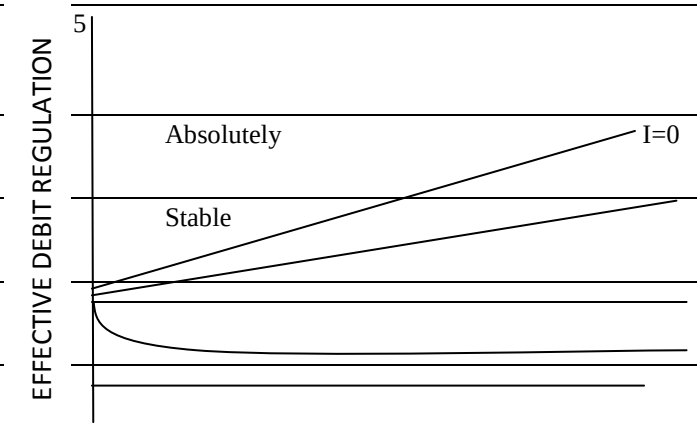



DELAYED DEBIT EFFECT

FIG6. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_6/dR_6 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.



-50 10

DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_6/dR_6 + N_6 dG_6/dR_6 + dB_6/dR_6)/M_6$). The immediate ($I_6$) and lagged ($L_6$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_6$(see 6b), is increased relative to the lagged part, $L_6$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_6 > I_6$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_6$is somewhat large than $I_6$ then the condition constraining prey regulation of recruitment ($dB_6/dR_6 < 0$) is

$$\frac{dD_6}{dR_6} + N_6\frac{dG_6}{dR_6} + \frac{dB_6}{dR_6} > \frac{(L_6 - I_6)M_6}{2} \qquad (10)$$

If $L_6 < I_6$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_6$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_6$

The - $D_6$ positions are defined by letting $\lambda_6 = iZ_6$ and separating (A1.2) into real and imaginary parts:

$$Z_6^2 + V_6 Z_6 \int_0^{\infty} \sin(Z_6 T_6) f(T_6) dT_6 - (X_6 + Y_6) = 0$$

$$Z_6 \left( U_6 - V_6 \int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6 \right) = 0$$

Noting that $X_6 + Y_6 = 0$ is we assume X+Y=0.It follows that $Z_6 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_6 = V_6 \int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6$$

If $U_6 > |V_6|$then given the constraints of $f(T_6)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A6.5)remains

greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

## ABSOLUTE STABILITY WITH A DELAY DUE TO CREDIT CONSOLIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_6^2 + (U_6 - V_6 + M_6)\lambda_6 - M_6\lambda_6 e^{-T_6\lambda_6} + [(U_6 - V_6)M_6 + Y_6] - [((U_6 - V_6 M_6 -]e^{-T_6\lambda_6} = 0$$

(For definitions see Appendix6) We follow the technique outlined in Appendix 6 and let $\lambda_6 = iZ_6$ and separate (A2.6) into real and imaginary parts

$$Z_6^2 + M_6 Z_6 \sin(Z_6 T_6) - (A_6 M_6 + Y_6) + (A_6 M_6 + X_6)\cos(Z_6 T_6) = 0$$

$$(A_6 M_6) Z_6 - M_6 Z_6 \cos(Z_6 T_6) + (A_6 M_6 - X_6)\sin(Z_6 T_6) = 0$$

Where $A_6 = U_6 - V_6$Using (A2.3), substitute for $Z_6^2$ and $Z_6$in (A2.2):

$$W_6^2(X_6^2 - 3M_6 A_6 X_6 - M_6^2(X_6 - Y_6) + 2M_6^3 A_6 + 2A_6^2 M_6^2]$$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_6 = 1 - \cos(Z_6 T_6)$.

In the absence of a delay the system is stable for $A_6 > 0,$ and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_6 > (X_6 - Y_6)/2\ M_6$ Thus a necessary condition for absolute stability is

$$A_6 > [\ X_6 - Y_6/2\ M_6$$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_6$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_6$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_6, A_6) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_6 \le 2$ *is* of interest First consider the extreme values

$$Q_6(0, A_6) = A_6^2\ (X_6 - Y_6)$$

$$Q_6(2, A_6) = (A_6 + 2\ M_6)^2\ (2A_6 M_6 - (X_6 - Y_6))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_6}{\partial A_6} = W_6^2\left(2M_6^3 + 4A_6 M_6^2 - 3M_6 X_6\right) + W_6\left(3M_6 A_6^2 - 2A_6 X_6 + 2M_6\left(2X_6 + Y_6\right)\right)$$

$$+2A_6\left(X_6+Y_6\right)$$

$=2M_6^2W_6^2(M_6+\mathrm{A}_6)+M_6A_6^2W_6+\left(2A_6M_6-X_6+Y_6\right)(A_6W_6+M_6W_6^2$

$+(2-W_6)((2M_6X_6+M_6Y_6)W_6+A_6(X_6+Y_6)),$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_6$quadratic in $W_6$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_6$has no minimum in the range $0 \le W_6 \le 2$.Defining

$Q_6\left(W_6,A_6\right)=a_6W_6^2+c_6 where\ a_6\ b_6\ and\ c_6$ Are functions of $A_6$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_6 \ge b_6 \ge 0$$

And

$$b_6^2 \le 4a_6c_6$$

$$Case(i): Y_6=0$$

$Let\ Y_6 = 0$ And try the substitution A= $X_6 = /2M_6\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_6 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_6 X_6 /$ $2M_6$) is positive for *W<2. Since* $Q_6$ increases with $A_6$(see A2.9), increasing $A_6$will generate no further solutions of $Q_6$=0. Hence the condition A>X/2M defines absolute stability when $Y_6 = 0$.

$Case\ (ii)$: $Y_6 = y_6\ X_6$ ,

$Let\ Y_6 = y_6\ X_6$And try the substitution $A_6 = X_6/(2 + A_6\ y_6)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_6/(2 + n_6 y_6)$ we have

$$W_6^2\ y_6\left[X_6 n_6 \frac{(1 + n_6 y_6)}{(2 + n_6 y_6)} + M_6^2\left(2 - n_6 + n_6 y_6\right)\right]$$

$$-W_6\ X_6\left[\frac{X_6(1 + n_6 y_6)}{M_6^2(2 + n_6 y_6)^2} + 2\ y_6\left(n_6 - 1\right) + \frac{X_6^2(6 + y_6)}{M_6^2(2 + n_6 y_6)} = 0.\ (A2.12)\right.$$

Since $n_6$ and $A_6$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_6$

Defining $E_6 = (2 + n_6 y_6), the\ condition (A2.66?) can\ be\ written\ as$

$[\frac{X_6(E_6 - 6)}{E_6^2}]^2[\frac{X_6^2}{M_6^2} - 4\ y_6\ E_6^2\frac{X_6}{M_6^2} + 4y_6(n_6 - 2)E_6^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in ($X_6 / W_6^2$) can be expressed as

$2E_6[y_6E_6+(y_6E_6(6+y_6)(E_6-n_6))^{6/2}]>\frac{X_6}{M_6^2}>2E_6[y_6E_6-y_6E_6(6+y_6)(E_6-n_6))^{6/2}]$

**ABSOLUTE STABILITY**

Hence if $X_6/M_6^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_6$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_6-6)}{E_6^2}\left(\frac{X_6}{M_6^2}\right)^2-2y_6\frac{(E_6+n_6^2y_6)}{E_6}\left(\frac{X_6}{M_6^2}\right)-4y_6(E_6-n_6)$$ 0

And has the solution

$\frac{X_6}{M_6^2}>\frac{E_6}{(E_6-6)}[y_6(E_6+n_6^2y_6)+[y_6^2(E_6+n_6^2y_6)^2+4y_6(E_6-6)(E_6-n_6)]^{6/2}]$

Given $E_6>n_6$ this lower limit to $X_6/M_6^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_6/M_6^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_6$ in the appropriate range and this limit is further lowered by (A2.64)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_6<n_6$ $(ie.,2A_6M_6^2<X_6-Y_6)$ then (A2.6)and

(A2.8)show that $Q_6$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_6}{M_6^2} > 2E_6[y_6E_6 - (y_6E_6(6 + y_6)(E_6 - n_6)]^{6/2}]$$

This equations throws light on the individual general ledger and Nature's general ledger as we call it : see various papers on the issue. powers, which are the energies of Matrka we have already discussed, are not the only obstacles the yogi must overcome. He must, for example, also resist the temptation to rest content with the miraculous yogic powers (siddhi) he acquires in the course of his spiritual develop- ment. Again to do this he must practice Yoga. Similarly, in order to pervade the Fourth State gradually through the other states in the manner proper to practice at the Individual level, the method is the same. He must practise the higher yoga of the Tantras which, turning his mind inwards and freeing it from discursive representations, allows him to penetrate into the Supreme Principle. 312 Once the yogi has attained this contemplative state, his main problem is to make it permanent. In the introverted state the gross external movement of the breath is suspended and with it the activity of the intellect, mind, individualised consciousness, powers of the senses and the ego. 313 When the yogi rises out of this state, he is liable to fall again into the lower order of creation generated by Maya if he does not maintain his awareness of the higher reality he has experienced and allows his awakened, illumined insight to be obscured by the dream-like vision of thought-constructs. 314 Naturally, the yogi must rise out of the introverted condition of suspension. It is inherent in the very nature of reality that it should move out of itself. 315 Pure, universal consciousness initially transforms itself into the vital breath 316 charged with the impression iyasana) of the power of awareness attained through introversion. By attending to the pulse (spanda) of the breath as it moves out of the absolute, the

yogi can develop an intuitive sense of the inherent unity of all he will perceive in the mental and physical spheres created by the outpouring of consciousness. In this way he realises that his own nature is everywhere present in all he perceives and that all things thus reside within him. Blessed with this insight his consciousness remains free and unlimited even at the individual level where the breath, mind, senses and body are active. If the yogi fails to do this, he finds himself once again beset by the strictures of his embodied existence and must, as before, try to pervade all his other states of consciousness with the aesthetic delight (rasa) and wonder (camatkara) of the Fourth State he experienced in contemplation. Again this means that he must strengthen his pure, empowered awareness that his universal nature manifests as all things. 317 In this way he discovers Siva's presence in every sphere of individualised consciousness ranging from the breath to outer objectivity. The yogi's mind then becomes tranquil and undistracted because wherever it may wander, the yogi perceives only Siva, his authentic nature. 318 Consciousness is thus freed of all external referents and the yogi's subjectivity is purified of all identification with the body or anything else that belongs to the objective sphere. The yogi then becomes detached from the opposites of pleasure and pain and is transcendentally free (kevalin). 319 yogis are again, however, liable to fall if he allows himself to get entangled in the play of opposites. This fall is more serious than the others because, although he is caught by the confining restrictions of individual- ised consciousness as before, he is now also affected by karma. Fleeing from pain in the pursuit of pleasure he is bound to act (karma) to minimise one and maximise the other and so is thrown down to the lowest level of embodied subjectivity (sakala). In order to regain his lost state, he must ascend gradually, by Siva's grace, from one order of subjectivity to the next and so free himself progressively of the limitations of the lower levels to gain the greater freedom and expansion of the higher. As he progresses, the objective sphere also evolves from the grossest perceptions of physical objects outside the lowest order to subjectivity, through to the subtler inner, mental

perceptions to finally reach the order of subjectivity that contains objectivity within itself and is free to externalise it at will. 320 The degree to which this process develops depends, as before, on the yogi's awareness of the Fourth State. In consonance with the general principle that the remedy should suit the defect, the yogi is instructed to seek this higher state of consciousness in the wonder (camatkara) or delight (ananda) he feels in moments of intense physical pleasure. At first he experiences this subtle consciousness for an instant in the subjective sphere. If he manages to catch hold of it, it becomes more intense as the cognitive and objective spheres are also gradually pervaded and vitalised by it. Occasions for this practice are, for example, the sense of satisfaction one feels after a good meal or the aesthetic delight one experiences when listening to good music or the pleasure of sexual union with the Tantric consort or even solitary sexual excitation. In these moments of delight the yogi can penetrate momentarily into his own authentic Siva-nature (sambhavavesa) through the empowered contact (saktasparsa) m he makes with it in the freedom of the pure subjectivity of the Fourth State. 322 If the yogi develops his awareness of this higher level of consciousness and maintains it, he eventually experiences it constantly. 323 Clearly, what prevents the yogi from attending to his state of consciousness rather than the circumstances which induce it is the craving for pleasure (abhila§a) born of ignorance — the source of every impurity which clouds consciousness. Craving directs the yogi's attention towards outer, worldly things and so he is caught in the net of thought- constructs. 324 To free himself of his worldly desires and reverse this binding extroversion, the yogi must eradicate its cause. To be freed of all the ups and downs of the path and no longer be tormented by the possibility of a fall, the yogi must see reality perfectly and completely. This insight is itself liberation and the moment it dawns the yogi is instantly freed. This sudden realisation is the goal of Tantric Yoga. E Tantra declares: "He who perceives reality directly, even for the brief moment it takes to blink, is liberated that very instant and never reborn again." 325 Although the yogi's body and mind

continue to function as before, they are like mere outer coverings 326 which contain, but do not obscure, the mighty, universal consciousness which operates through them. The yogi's body is (=) the universe, the senses (=) the energies that (eb, eb+)vitalise it, his mind (=) Mantra, the rhythm of his breath (eb) the pulse of time and his inner nature pure, dynamic consciousness. Raised above all practice, and hence all possibility of falling to lower levels, the yogi realises that he has always been free 327 and that his journey through (e&eb) the **dark land of Maya was nothing but a dream, a construct of his own imagination** The channels of the senses by which the yogi enjoys the pleasure of the objects of sense are those through which he fills the three worlds with the consciousness emitted from his own Heart. 78 **The Doctrine of Vibration T**he senses are outer, subsidiary cycles of conscious energy (anucakra), revolving around the core or primary cycle (mukhyacakra) of pure consciousness. They scintillate with the diversity of sensations and resonate in the field of their awareness like the feathers of a peacock unfurled and folded back in the ecstasy of its dance. 79 Ultimately it is Siva Himself Who, for His own pleasure, 80 extends and withdraws them, 81 not the individual soul. The outward movement of the senses is the creative flow of consciousness from the subject who, like an ocean, is the source of the waves of their activity and the final resting place of the flow of the rivers of sensations and thought-forms that unfold through the arising of the power of each of the senses. 82 The Stanzas on Vibration declare: Indeed the individual soul does not activate the impulse of the will [which directs the body's activity] by himself alone but through his contact with [his] own [inner] strength {bald) made in such a way that he identifies with it [thus acquiring its power]. 83 The close contact between the senses and universal consciousness vitalises them and sets them in motion. They share in the properties of consciousness, just as a heated iron ball can burn, cook or give out light. 84 Spanda doctrines does not agree with the Sarnkhya view that the individual soul (purusa) alone stimulates activity in the objective sphere (prakrti) by the contact he makes with it. The soul can do

this only because universal consciousness animates him. Moreover, consciousness, which has the power to impel this activity, cannot itself be inactive as the Sarnkhya maintains. The influence it has on lifeless matter implies an active exertion. 85 This activity is the subtle, inner outpouring of the Self (svatmocchalatta) within the Self. Outside the confines of time and space, in the centre of the circle of the senses, this inner activity initiates the expansion (prasara) of the flow of conscious energy. As the pure exertion (udyoga, udyama) of Being (bhava), it is the source of the initial unfolding (unmesa) of Becoming. 86 Ksemaraja explains: [By] 'exertion' (udyama) is [meant] the emergence of the supreme [level of] intuition (pratibha) which is the sudden outpouring of the [immanent and transcendental] unfolding of the reflective awareness of pure consciousness. 87 The direct, intuition (pratibha) of reality, free of thought-constructs, is the supreme power (parasakti) of consciousness. The yogi can gain Divine **Body and Sacred Circle of the Senses 153** this pure intuitive awareness by grasping the initial expansion (unmesa) of consciousness which engenders the emergence of thought and the activity of the senses. It can be discovered in the subtle Centre, between one thought and the next. The Stanzas on Vibration teach: The expansion of consciousness that takes place when one is engaged in a single thought should be known to be the source from whence another arises. One should experience that for oneself! 88 This does not mean that the unfolding expansion of consciousness {unmesa) is an intermittent, transitory occurrence. 89 As the source of all thought, including the notions of past, present and future, 90 it stands outside the confines of time. Consciousness is in a state of perpetual expansion. It is an ever developing process, a tiny part of which we observe as the flow of events around us and the thoughts and feelings within us. The totality of this process can never be grasped objectively. Objectivity is a state of fragmentation (bheda) in which discrete elements are discernible and mutually distinguished as individual, specifically definable entities. It is the realm within which thought and language operate. To grasp reality in its completeness, we must go

beyond the partial representations of thought and speech. To experience the primordial source and basis of all things, we must pierce through the outer periphery of thought and plunge into the Centre, to discover the instant in which thought, and with it, the sense of diversity, initially emerges. Ksemaraja cites a practice from the Vijnanabhairava through which the yogi can achieve this state of awareness: Checking [the movement of] attention (cit) once it has quit its object, it must not [be allowed] to move to any other; contemplation (bhavana) then blossoms forth by [experiencing] the state in the centre between them (tanmadhyabhava). 91 Left unattended and undisciplined, awareness spontaneously shifts from one object of thought to (e&eb) the next. The transition from one to the other entails a movement through (e&eb) a state of pure indeterminate (nirvikalpa) awareness divested of all thought-forms and perceptions. The yogi who succeeds in checking the movement of his attention, experiences within himself the subsidence of thought into the emptiness of the Centre which unfolds full of the powerful pulsation of consciousness within which all objectivity merges and becomes one 92 in the trans- cendental outpouring (lokottarollasa) of its aesthetic rapture. 93 turning in on himself, the yogi must exert his power of awareness to (e) discern the vitality of (e) his own blissful Spanda nature. Just as a distant object becomes (=) more clearly visible by (e) exerting an effort to (e) perceive it, in the same way, through (e&eb) practice, the yogi's own Spanda nature becomes (=) progressively more evident to him, 94 The effort the yogi exerts to free himself of ignorance is essentially one with the force of awareness that impels the senses and mind. This inner exertion is the basis of every means to liberation. Ultimately bestowing the bliss of self- realisation, it carries the yogi along the path to salvation with all its twists and turns like wind which blows a cloth in various directions. 95 As the basis and ultimate goal of all yogic practice, this exertion is the pure wonder (camatkdra) inspired by the realisation of the essential identity of consciousness (cit) and being (sat). 96 As such it is the divine radiance (sphuratta) of the Heart of one's own consciousness 97 where the powers of

will, knowledge and action merge in the harmony (samarasya) of bliss. 98 Thus the highest form of worship (varivasya) is the contemplation of the inner strength (bald) or universal exertion of one's own true nature. 99 "One should examine [that reality]," says Ksemaraja, "with the exertion which is Bhairava's nature and is true service (sevana) to the integral inner nature [of all things] and the unfolding of the activity of one's own vitality (ojas) which withdraws all duality [back into consciousness]." 100 Mindful of the eternal joy (ahlada) of his inner strength, the yogi must give up all desire for the fleeting moments of petty pleasure (ksanikasukhalava) ]0] he may glean from outer objects. He must even give up the desire to achieve enlightenment and a reality (tattva) which can never be realised through a personal act of will. 102 In short, he must die to himself in an act of profound faith (sraddha) and adoration (adara) to find rest in the supreme reality free from the binding dichotomy of means and objectives (upeyopayabhava) — the reality of his perfected effort. The yogi, now full of devotion for this reality unfolding within himself and the universe (antarmukhatattva), merges into it and, by achieving the pure consciousness and bliss (cidananda) of the power (bala) of his own nature, achieves oneness. 103 The yogi is firmly established in his true nature by recognising the creative unfolding (unmesa) of consciousness through which he is able to know all things, whether in the present, past or future. 104 Master of the vitality which gives life to his mind and body, he can strengthenit directly without need of food or drink, 105 or quit it at will to vivify Divine Body and Sacred Circle of the Senses 155 that of another. 106 Experiencing the eternal satisfaction of his own uncreated nature, 107 the yogi is freed of disease and the infirmities of old age. 108 Inspired by this vitality, he has the inner strength to overcome the severest hardships and obstacles in the way of his ultimate goal. 109 The Stanzas on Vibration explain: Lassitude ravages the body and this arises from ignorance but if this is eliminated by an expansion of consciousness (unmesa), how can [ignorance], deprived of its cause, continue to exist? 110 Doubt (sanka) is the source of every spiritual ailment. 111 A

man in doubt about his true nature, and the path to follow by which he may come to realise it, is constantly overcome by difficulties. 112 "Due to ignorance" teaches the Sarvacaratantra, "the deluded man is in doubt and thus [suffers the cycle] of birth and death." 113 Caught between conflicting alternatives, the mind shifts from one to the other becoming more entangled in its own thoughts as it does so. Doubt thus contracts consciousness 114 sullying it with the turbulence of confused thinking. 115 At once both the root of "the ancient tree of transmigration" and the first sprout of its seed, 116 doubts deprives man of the innate bliss (sahajananda) of his own nature. 117 Like a thief it steals away the wealth of true knowledge and reduces man to an imagined state of poverty in which he feels hemmed in on all sides by constraints and limitations. 118 Ignorance and its sister, doubt, are the essence of all impurities (mala) which sully consciousness. To the degree in which the yogi is cured of this ailment, namely, the lassitude (glani) of doubt, his true nature becomes manifest, just as raw gold when heated is freed from dross. When the yogi recognises the all-powerful expansion (unmesa) of his consciousness, he is projected beyond the realm of relative distinctions between virtue and vice. Thus overcoming all doubt as to what he should or should not do, he penetrates into the pulsing Heart of Bhairava's consciousness. 119 As the Stanzas on Vibration say: An individual who [though] desirous of doing various things [but] is incapable of doing them due to his innate impurity, [experiences] the Supreme state when the disruption (ksobha) [of his ego] ceases. Then [the soul realises] that his [true] uncreated nature is [universal] agency and perceiving subjectivity and so he knows and does whatever [he] desires. 120 the individual soul is(=) a combination of (e) conscious (cit) and 156 The Doctrine of Vibration unconscious (acii) elements 121 brought together when he allows the power of his authentic nature to be obscured by (e) his own impurity. This impurity is (=) three-fold. The first is technically called (=) the 'Impurity of Individuality' (anavamald). Due to (e) this impurity, the individual soul fails to (e) recognise his all-embracing fullness (purnata) and is disturbed by (e) his craving

for (e) experience (bhoga) based on (e) the unreasoned assumption that he is (=) incomplete (apurnd) and in need of (e) something outside himself. 122 Thus it manifests as (=) desire or attachment (rdga). ]23 At the same time it is (=) primordial ignorance, 'the cause of the sprout of (e) transmigration'. 124 As such, it works in two ways. Firstly, it deprives (e) individual consciousness of freedom. The yogi affected by (e&eb) this impurity is incapable of (e) entering or emerging out of the states of contemplation as he wishes. Secondly, this impurity renders (eb) the individual unconscious of his freedom, which for(e) the yogi means that he is (=) liable to lose consciousness of (e) his true nature when he is (=) not in a state of (e) contemplative absorption. 125 The second impurity is (=) called the 'Impurity of Maya' (maylyamald). This comes into operation when consciousness has been contracted by (e) the Impurity of Individuality. It limits (e) the power of knowledge operating in (e&eb) the individual soul through (e&eb) the inner and outer organs of (e) perception. 126 Due to this impurity he perceives (e) diversity everywhere. 127 Finally, the third impurity, that of Karma, comes into play when (e) the individual, deprived of (e) the freedom and knowledge of (e) unconditioned consciousness, acts (e&eb) in his limited way prompted by (e) desires and fears for (e) his personal gain. Disturbed by (e) these three impurities the soul cannot find rest in (eb) himself.

And

$$E_6 > n_6$$

Replacing $E_6$ and $n_6$ in these conditions leads to the form

If $A_6^2 > 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6$.

If $A_6^2 < 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6 .+ (Y_6 - A_6^2/2)^2/2M_6(X_6 - Y_6)$.

Note that by normalizing $A_6$ by $M_6$ and $X_6$ and $Y_6$ by $W_6^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_6^2 + (M_6 + U_6)\lambda_6 - (M_6 + V_6)\lambda_6 e^{2T_6\lambda_6} + (M_6 U_6 + Y_6)$$

$$-[M_6(U_6 + V_6) - X_6]\, e^{2T_6\lambda_6} + M_6 V_6\, e^{2T_6\lambda_6} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_6 = \mathrm{i}\, Z_6$ *and* separate

(A3.1)into real and imaginary parts

$$Z_6^2 + (M_6 + V_6) Z_6 \sin(Z_6 T_6) - (U_6 M_6 + Y_6) + (M_6(U_6 + V_6) - X_6)$$

$$\mathrm{X}\cos(Z_6 T_6) - 2M_6 V_6 \cos^2(Z_6 T_6) + M_6 V_6 = 0$$

$$(M_6 + U_6) Z_6 - (M_6 + V_6) Z_6 \cos(Z_6 T_6) + (M_6(U_6 + V_6) - X)$$

$$\mathrm{X}\sin((Z_6 T_6) - M_6 V_6 \sin(Z_6 T_6) = 0$$

Using (A.3.3) to eliminate $Z_6$ and $Z_6^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_6 =$

$1 - \cos Z_6 T_6$ and $A_6 = U_6 - V_6$ we obtain

$$4M_6^2 V_6^2 W_6^4 + 2M_6 V_6 W_6^3[(M_6 - V_6)^2 + A_6(M_6 + V_6) + 2(A_6 M_6 - X_6)]$$

$$+ W_6^2 [M_6 - V_6)^2 (2 A_6 M_6 + Y_6 - X_6) + (2A_6 M_6 - X_6)^2 - A_6 (M_6 - V_6)(2A_6 M_6 - X_6)$$

$$+ M_6 V_6 (A_6^2 + 4(X_6 + Y_6))]$$

$$+ W_6 [A_6^2 + 2X_6)(2 A_6 M_6 - X_6) + 2A_6 Y_6 (M_6 + V_6) - M_6 A_6^3] + A_6^2 (X_6 M_6 - Y_6) = 0$$

Being a quartic in $W_6$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_6 > 0$$

$$and$$

$$C_6 > (X_6 - Y_6)/2 M_6$$

$recalling\ that\ A_6 = U_6$ and defining $C_6 = U_6 + V_6$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_6 (0, U_6 V_6 = A_6^2 (X_6 + Y_6)$$

$$Q_6 (2, U_6 V_6) = (C_6 + 2M_6^2)^2 (2C_6 M_6 + Y_6 - X_6)$$

So that given (A3.6) the extreme $W_6$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_6$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_6$ the existence of a D-position at this point is

determined by letting cos $(Z_6 T_6) = -1\ in\ (A3.3)$

ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 e^{-\tau_6\lambda_6} + X_6 + Y_6 = 0 \qquad (A1.1)$$

Where $V_6 = dB_6/dR_6$ $U_6 = dD_6/dR_6 + N_6 dG_6/dR_6$ $X_6 = N_6G_6dF_6/dR_6$ and $Y_6 = -N_6G_6dM_6/dR_6$ (evaluated at the equilibrium). Note that $U_6X_6$ and $Y_6$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_6$ as $(T_6)$ . (It follows that $f(T_6)$ 0 for all $T$, and $\int_0^\infty f(T_6)dt = 1.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 \int_0^\infty e^{-\tau_6\lambda_6} f(T_6) + X_6 + Y_6 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MODULE NUMBERED TWO**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed

the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).. Haimovici's and Nunny's models have been applied to some of the Deleuzean and Heideggerian problems too : Scoring on that count are :

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time

delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\left[F\left(R(t)\right) - M\left(R(t)\right)\right]$$

(3)

$$\frac{dR(t)}{dt} = B\left(R(t-T)\right) - D\left(R(t)\right) - N(t)G\left(R(t)\right).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t - T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if

only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

TIME DELAY

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

## MATURATION AND CONSUMMATION OF CONFEDERATION OF DEBIT AND CREDIT TYPES:

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t-T_7)F_7(R_7(t-T_7) - N_7(t)M_7((R_7(t) \qquad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7} + \frac{N_7 dG_7}{dR_7} - \frac{dB_7}{dR_7} > \frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7} + \frac{dM_7}{dR_7}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes $(L_7)$ and the other derived from the non-delayed or "immediate" processes$(I_7)$:as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_7 = \frac{N_7 G_7}{2M_7}\frac{d(7nF_7)}{dR_7} (\geq 0) \qquad (7a)$$

$$I_7 = -\frac{N_7 G_7}{2M_7}\frac{d(7nM_7}{dR_7} \ (\geq 0) \qquad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7 > L_7 - I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_7 > L_7 - I_7, \quad \text{when} \quad P_7^2/4 > I \qquad (8a)$$

$$P_7 > L_7 - I_7 + \frac{(I_7 - P_7^2}{(L_7 - I)} \quad \text{when} \quad P_7^2/4 < I. \qquad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

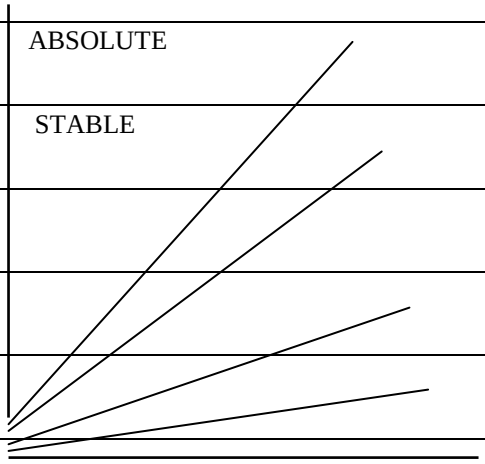



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit

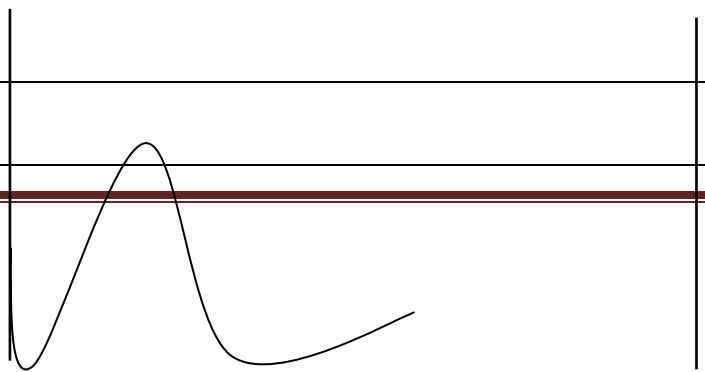
(like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

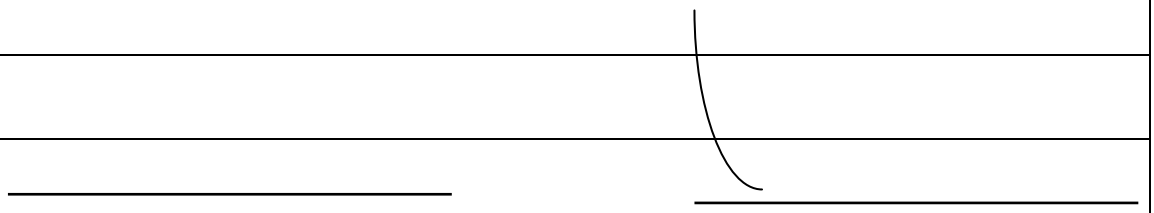
It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the





TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t-T_7)F_7\big(R_7(t-T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t-T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation $(dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7)$ and the effective prey regulation $(dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7)$ must be positive.

The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the

total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

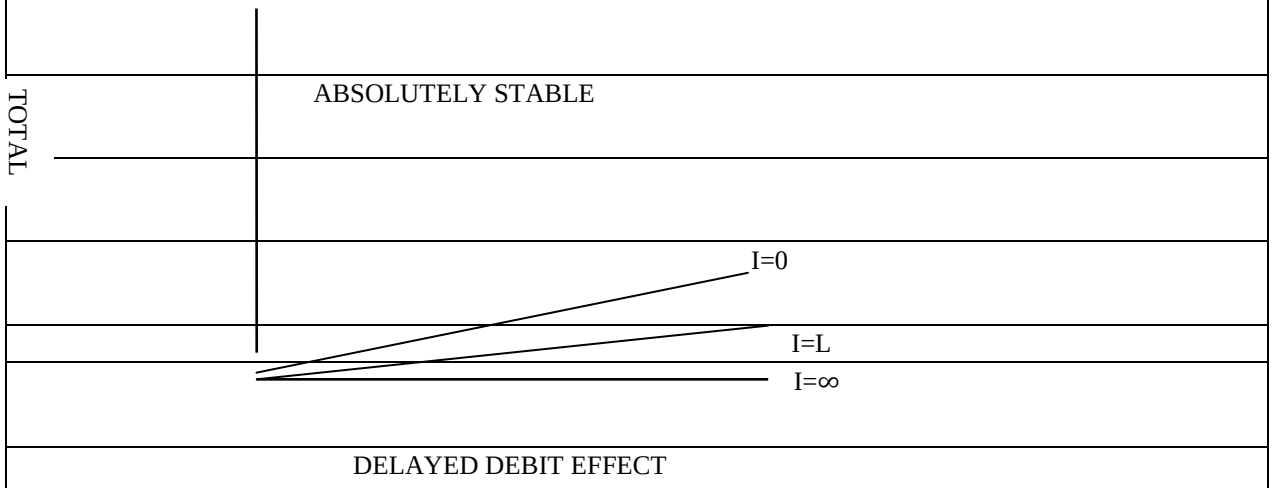



FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

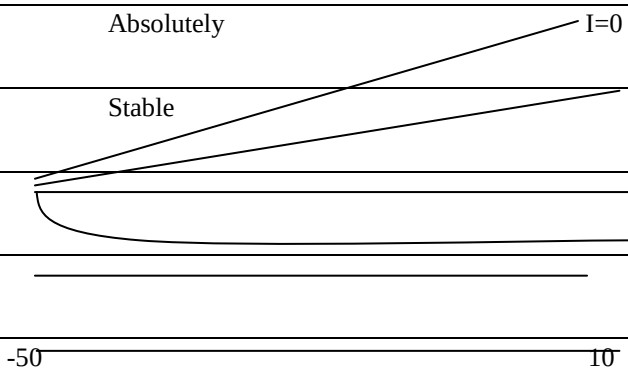
Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.

5

DEBIT REGULATION




**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay.

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^\infty \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7\left(U_7 - V_7\int_0^\infty cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7\right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7\int_0^\infty cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$then given the constraints of $f(T_7)$, $(A1.5)$*can never be satisfied Co*nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7\lambda_7 e^{-T_7\lambda_7} + [(U_7 - V_7)M_7 + Y_7] - [((U_7 -$$

V 7M 7−]$e^{-T7\lambda 7}$=0

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$Z_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7)\cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7)\sin(Z_7 T_7) = 0$$

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $Z_7^2$ and $Z_7$in (A2.2):

$W_7^2 (X_7^2 - 3 M_7 A_7 X_7 - M_7^2 (X_7 - Y_7) + 2 M_7^3 A_7 + 2 A_7^2 M_7^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_7$=1-cos $(Z_7 T_7)$.

In the absence of a delay the system is stable for $A_7 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7 > (X_7 - Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$$A_7 > [\ X_7 - Y_7/2\ M_7$$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_7, A_7) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_7 \le 2$ *is* of interest First consider the extreme values

$$Q_7(0, A_7) = A_7^2\ (X_7 - Y_7)$$

$$Q_7(2, A_7) = (A_7 + 2\,M_7)^2\,(2A_7 M_7 - (X_7 - Y_7))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7 M_7^2 - 3M_7 X_7\right) + W_7\left(3M_7 A_7^2 - 2A_7 X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7+Y_7\right)$$

$=2M_7^2W_7^2(M_7+\mathrm{A}_7)+M_7A_7^2W_7+\left(2A_7M_7-X_7+Y_7\right)(A_7W_7+M_7W_7^2$

$+(2-W_7)((2M_7X_7+M_7Y_7)W_7+A_7(X_7+Y_7)),$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $W_7$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$has no minimum in the range $0\leq W_7\leq 2$.Defining

$Q_7\left(W_7,A_7\right)=a_7W_7^2+c_7 where\ a_7\ b_7\ and\ c_7$ Are functions of $A_7$we can define the condition for a minimum in the range 0-2

condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7\geq b_7\geq 0$$

And

$$b_7^2\leq 4a_7c_7$$

$$Case(i):Y_7=0$$

$Let\ Y_7=0$ And try the substitution A= $X_7=/2M_7\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_7 X_7 /$ $2M_7$) is positive for *W<2. Since* $Q_7$ increases with $A_7$(see A2.9), increasing $A_7$will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

$Case\ (ii): Y_7 = y_7\ X_7$ ,

*Let* $Y_7 = y_7\ X_7$And try the substitution $A_7 = X_7/(2 + A_7\ y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + n_7 y_7)$ we have

$$W_7^2\ y_7\left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2\left(2 - n_7 + n_7 y_7\right)\right]$$

$$- W_7\ X_7\left[\frac{X_7(1 + n_7 y_7)}{M_7^2(2 + n_7 y_7)^2} + 2\ y_7\left(n_7 - 1\right) + \frac{X_7^2(7 + y_7)}{M_7^2(2 + n_7 y_7)}\right] = 0.\ (A2.12)$$

Since $n_7$ and $A_7$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7), the\ condition (A2.77?) can\ be\ written\ as$

$[\frac{X_7(E_7 - 7)}{E_7^2}]^2[\frac{X_7^2}{M_7^2} - 4\ y_7\ E_7^2\frac{X_7}{M_7^2} + 4y_7(n_7 - 2)E_7^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$2\ E_7[y_7 E_7 + (y_7\ E_7(7 + y_7)(E_7 - n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - y_7 E_7(7 + y_7)(E_7 - n_7))^{7/2}]$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0⁄$W_7$ ⁄2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7 - 7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2 - 2y_7\frac{(E_7 + n_7^2 y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right) - 4y_7(E_7 - n_7) \text{⁄0}$$

And has the solution

$$\frac{X_7}{M_7^2} > \frac{E_7}{(E_7 - 7)}[y_7(E_7 + n_7^2 y_7) + [y_7^2 (E_7 + n_7^2 y_7)^2 + 4y_7(E_7 - 7)(E_7 - n_7)]^{7/2}]$$

Given $E_7 > n_7$this lower limit to $X_7 / M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7 / M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$in the appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7 < n_7$ $(ie., 2A_7 M_7^2 < X_7 - Y_7$)then (A2.7)and (A2.8)show that $Q_7$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - (y_7 E_7(7 + y_7)(E_7 - n_7)]^{7/2}]$$

And

$$E_7 > n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2 > 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.

If $A_7^2 < 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7 + (Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$ and $X_7$ and $Y_7$ by $W_7^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSILIDATAN**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7\lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]e^{2T_7\lambda_7} + M_7 V_7 e^{2T_7\lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i}Z_7$ *and* separate

(A3.1)into real and imaginary parts

$$Z_7^2 + (M_7 + V_7)Z_7 \sin(Z_7 T_7) - (U_7 M_7 + Y_7) + (M_7(U_7 + V_7) - X_7)$$

$$\mathrm{X}\cos(Z_7 T_7) - 2M_7 V_7 \cos^2(Z_7 T_7) + M_7 V_7 = 0$$

$$(M_7 + U_7)Z_7 - (M_7 + V_7)Z_7 \cos(Z_7 T_7) + (M_7(U_7 + V_7) - X)$$

$$X \sin\left(\left(Z_7 T_7\right) - M_7 V_7 \sin\left(Z_7 T_7\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_7$ and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - cos\ Z_7\ T_7$ and $A_7 = U_7 - V_7$ we obtain

$$4 M_7^2 V_7^2 W_7^4 + 2 M_7 V_7\ W_7^3 [(M_7 - V_7)^2 + A_7 (M_7 + V_7) + 2(A_7 M_7 - X_7)]$$

$$+ W_7^2 [M_7 - V_7)^2 (2 A_7\ M_7 + Y_7 - X_7) + (2A_7\ M_7 - X_7)^2 - A_7 (M_7 - V_7)(2A_7\ M_7 - X_7)$$

$$+ M_7 V_7 (A_7^2 + 4(X_7 + Y_7))]$$

$$+ W_7 [A_7^2 + 2X_7)(2 A_7\ M_7 - X_7) + 2A_7 Y_7 (M_7 + V_7) - M_7 A_7^3] + A_7^2 (X_7\ M_7 - Y_7) = 0$$

Being a quartic in $W_7$ (A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7 > 0$$

$$and$$

$$C_7 > (X_7 - Y_7)/2\ M_7$$

$recalling\ that\ A_7 = U_7$ and defining $C_7 = U_7 + V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0, U_7 V_7 = A_7^2 (X_7 + Y_7)$$

$$Q_7\left(2, U_7 V_7\right) = \left(C_7 + 2M_7^2\right)^2 (2C_7 M_7 + Y_7 - X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1\ in\ (A3.3)$

## ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

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# Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price, University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss**

(777) In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them.

(778) Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors.

**Classical vs. relativistic black holes**

(779) Something like a black hole exists within (eb) Newton's classical theory of gravity.

(780) In that theory, an energy argument tells us that there is (=) an escape velocity vesc=2GM/R−−−−−−−√ from (e) the surface of any spherical object of (e) mass M and radius R

(781) If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity.

(782) Thus the condition for such an "unseeable" object is R<2GM/c2

(783) In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape.

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| (784) This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity. |
| (785) To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it. |
| (786) The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c . |
| (787) This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed. |
| (788) An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity. |
| (789) Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable. |

**Stationary black holes**

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| (790) In the Newtonian theory, gravity is described by (e) the potential $\Phi$. |
| (791) Inside a spherical object the form of (e) $\Phi(r)$ depends on (=) the interior structure, but in the vacuum outside (e) matter the potential $-GM/r$ depends only on (=) the interior mass. |
| (792) Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object. |
| (793) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward. |
| (794) This horizon radius $RH=2GM/c2$ is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory. |
| (795) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution. |
| (796) Radius RHhere actually means that the area of the event horizon is (=) $4\pi R2H$. |
| (797) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution. |
| (798) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity. |
| (799) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own. |
| (800) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole. |
| (801) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density. |
| (802) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed. |
| (803) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole. |

(804) In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter.

(805) Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum.

(806) In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior.

(807) The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object.

(808) Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum.

(809) These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.]

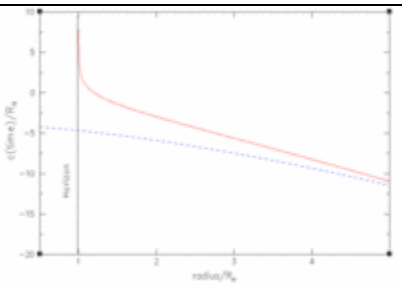
(810) All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time.

(811) But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?"

(812) The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point.

(813) One of these coordinates is (=) called the "coordinate time."

(814) Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate.



(815) Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole

(816) At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole.

(817) Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations.

(818) For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree

(819) An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events.

(820) The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving.

(821) Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle.

(822) The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return.

(823) But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon.

(824) Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future.

(825) The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?"

**Black hole parameters**

(826) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin).

(827) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun.

(828) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits.

(829) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form.

(830) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds.

(831) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2.

(832) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed.

(833) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer.

(834) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2)

(835) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials.

(836) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable

(837) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite.

(838) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole.

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| (839) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars. |
| (840) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity. |
| (841) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity. |
| (842) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite. |
| (843) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity. |
| (844) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum. |
| (845) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship. |
| (846) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= $(4\pi G2/c4)[(M+M2-a2\sqrt{\phantom{xx}})2+a2]$. |
| (847) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored. |

**Dynamical black holes**

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| (848) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape. |
| (849) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, **a horizon can be born**. |
| (850) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid. |
| (851) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward. |
| (852) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations. |
| (853) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes. |
| (854) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3. |
| (855) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/ |
| (856) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers. |
| (857) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission. |

(858) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.

(859) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.

(860) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.

(861) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.

(862) The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature $6\times10-8$ (Mo/M) K.

(863) Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.

**Astrophysical black holes**

(864) Black holes in our Universe can be grouped as: primordial black holes, stellar-mass black holes, and supermassive black holes. The first is highly speculative; the second and third are broadly accepted.

(865) Primordial black holes of all masses are postulated to have formed from (e) quantum fluctuations in the early Universe, but those with (e&eb) mass less than around 1013kg would have already evaporated due to Hawking radiation.

(866) No significant observational evidence has yet been found of (e) the existence of these objects [MacGibbon and Carr (1991)].

(867) Stellar-mass black holes, ranging from a few to a few tens of solar masses, are (=) normal but rare endpoints in (eb) the evolution of massive stars.

(868) When a star exhausts its nuclear fuel and cools, it must collapse unless (e) it is supported by nonthermal forces.

(869) It is known that such nonthermal forces cannot (e) resist gravitational compression for (e) masses greater than around 1.5 to 3Mo. (The uncertainty is due to (e) uncertainty in our knowledge of nuclear physics at high densities.)

(870) There are many stars far more massive than this, but in their death throes massive stars expel (eb) a great deal of their mass in supernova explosions

(871) Even stars initially as massive as 20-30Mo may blow off enough mass to (e) leave a remnant neutron star smaller than 1.5Mo.

(872) Stars more massive than 20-30Mo may form (eb) a neutron star core, only to have it collapse due to (e) fallback of material from (e) the stellar mantle. Still more massive stars (above ~40Mo) may form (eb) a black hole directly in collapse, with or without (e&eb) a supernova explosion.

(873) These masses are (=) rather uncertain at present, as our understanding of supernova explosion mechanisms is (=) still evolving [Woosley and Janka (2005), Muno (2006)].

(874) Collapsing gas clouds larger than ~100Mo are expected to dissipate from (e) radiation pressure, which would prevent (e) more massive stars from forming, and consequently set (eb) an approximate upper limit for stellar mass black holes.

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| (875) Observational evidence for stellar-mass black holes comes primarily from (e) X-ray astronomy. |
| (876) A black hole in a close binary orbit can form (eb) an accretion disk (see Figure 1) of matter pulled off (er) a normal stellar companion. In its inspiral, disk material is heated by shearing, and becomes (=) a strong X-ray emitter. |
| (877) If observations reveal a point-like X-ray emitter with (e&eb) a mass inferred from (e) orbital dynamics to be (=) above that possible for a neutron star, it becomes (=) a black hole candidate. |
| (878) t present there are over a dozen such black hole candidates, including the first object to be identified as (=) a black hole, the X-ray source Cygnus X-1 (estimated mass 10±3Mo) [Casares (2006)]. |
| (879) The other class of observationally-supported black holes is (=) that of supermassive black holes, ranging from hundreds to (e&eb) billions of solar masses. |
| (880) Evidence for the existence of black holes at (eb) the upper end of this range is overwhelming. |
| (881) By contrast, the existence of black holes with (e&*eb) mass roughly of order 102 to 104Mo, referred to as "intermediate-mass black holes," is speculative. |
| (882) The supermassive holes may have begun as (=) primordial or stellar-mass black holes, but have grown through (e&eb) absorption of stars or gas or through mergers with (e&eb) other holes [Ferrarese and Ford (2005)]. |
| Figure 4: Observed orbits of stars near (eb) the dark compact object at the center of (e) our Milky Way galaxy: still frame from an animation. Animation created by Prof. Andrea Ghez and her research team at UCLA, from data sets obtained with the W. M. Keck Telescopes: http://www.astro.ucla.edu/~ghezgroup/gc/ |
| Evidence for supermassive black holes originally came from (e) quasars, small but intense radio sources seen at (eb) cosmological distances. |
| Their huge luminosities implied (eb) high mass, while (e) their rapid variability implied (eb) extremely small size. |
| More recently, measurements of Doppler shifts of stars and gas in the centers of galaxies have shown (eb) that compact objects of mass greater than 106Mo reside in (eb) the cores of most galaxies; the very small size of these central objects rules out any plausible alternative to the black hole explanation. |
| The black hole at the center of our own Milky Way galaxy has (e) a measured mass of ~3×106Mo , based on (e) the orbits of stars near the radio source Sagittarius A* associated with (e&eb) the Galactic core [Ghez et al. (2005); see Fig. 4]. |
| (888) Black holes are of particular interest for (e) gravitational waves, and vice versa. |
| (889) For gravitationally-bound systems, the typical maximum gravitational wave amplitude (dimensionless strain) is (=) h~(2GM/c2)2/(RD)~R2H/(RD) , where D is the distance to the system, M and R are the mass and size of (e) the system, and RH is the horizon radius of (e) a black hole of the same mass. |
| (890) From this it is clear that strongest gravitational waves will involve (e&eb) systems at or near black hole compactness: in particular, systems containing (er) black holes in close proximity to one another [Flanagan and Hughes (2005)]. |
| (891) Conversely, a system containing only black holes would not be expected to radiate (eb) any sort of radiation other than gravitational waves, and in any system containing (e) black holes, gravitational waves provide (eb) the most accurate probes of conditions deep in (eb) their gravitational wells. |
| (892) In particular, gravitational wave observations could definitively test (e&eb) whether a compact mass is truly a black hole [Cutler and Thorne (2002)]. |

(893) The strongest intrinsic source of gravitational waves in the present Universe is (=) the collision of two comparable-massed black holes.

(894) A sufficiently tight black hole binary system will lose (e) energy through emission of (e&eb) gravitational waves, causing (eb) the orbit to circularize and then slowly shrink over time.

(895) Figure 5: Gravitational waves from (e) a black hole binary inspiral: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/

(896) When the orbital radius approaches (eb) the Schwarzschild radius of the system, complex nonlinear dynamics come into play (e&eb) and the final stages of (e) the inspiral and merger can only be modeled through (e&eb) numerical simulations.

(897) These simulations are quite challenging, but recent breakthroughs [Pretorius (2005)(See Fig. 5), Campanelli et al.(2006), Baker et al.(2006)] have led to (eb) the first complete inspiral-merger gravitational waveforms.

(898) After the black holes have coalesced, they form (eb) a single highly-distorted black hole which quickly settles down to (e) a quiescent Kerr state through (e&eb) mission of quasinormal "ringdown" gravitational waves [Pretorius (2005), Campanelli et al. (2006), Baker et al. (2006)].

(899) Another gravitational emission mechanism for supermassive black holes would come from (e) their capture of compact stellar objects, particularly stellar-mass black holes.

(900) Unlike comparable-mass mergers, these captures would be on (eb) highly eccentric orbits that would not circularize before merger, resulting in (eb) intricate rapidly precessing orbits in (eb) the final months or years of inspiral.

**(901)** The consequent gravitational-wave signal would allow (eb) for exquisitely precise measurements of the parameters of (e) the supermassive black hole, as well as tests of any deviation from (e) the predictions of general relativity [Cutler and Thorne (2002)].

**(902)** For stellar-mass black holes (tens of Mo), the final stages of inspiral and merger occur at (eb) frequencies of hundreds to thousands of Hz, a regime targeted by (e) ground-based gravitational-wave observatories.

**(903)** Systems of supermassive black holes that are (=) millions of Mo emit at frequencies of tens of (e) mHz, which will be detectable by (e) proposed space-based (e) gravitational-wave detectors.

**(904)** Black holes in the billions of Mo, or a stochastic background of signals from (e) more moderate supermassive black holes at (eb) earlier stages of their inspiral, might emit (eb) waves at nHz frequencies that would be measurable through (e&eb) years-long correlated timing measurements of radio pulsars [Jenet et al. (2005)].

**An equation to quantify the origins of life on other planets July 5, 2016 by Bob Yirka report**

(905) (Phys.org)—A pair of researchers, one with the Columbia Astrobiology Center in New York, the other with the University of Glasgow in the U.K. has come up with a mathematical equation that when (e) solved is meant to offer (eb) a means for estimating how often life begins on other planets. In their paper published in Proceedings of the National Academy of Sciences, Caleb Scharf and Leroy Cronin describe their equation, how they came up with it and why they believe it might become more useful as scientists learn more about the true nature of other planets and solar systems.

(906) Back in the 1960's, an astronomer named Frank Drake came up with a formula for (e) estimating the likely number of (e) alien civilizations that might be capable of transmitting radio signals in such a way as to be (=) recognizable by receivers here on earth.

(907) That equation, named quite naturally the **Drake Equation** has been the unofficial standard-bearer for half a century, despite it having no solution because (e) some of the parameters are still unknown.

(908) The same holds true for the new equation developed by (e) Scharf and Cronin, but the results given should the unknowns ever be discovered would offer a different sort of answer than Drake was looking for.

(909) To build their formula, the duo came up with several parameters to solve for (e) **abiogenesis**, which is the likelihood of an event occurring that leads to (eb) life beginning, they included (e) the number of possible building blocks, the mean number of such blocks per possible organism, the availability of building blocks that might exist during a given time period, expressed as a fraction and the probability that the existence of the building blocks would actually lead to (eb) life starting in a given unit of time. It looks like this: **Nabiogenesis (t)= Nb · 1/no · fc · Pa · t**

(910) In essence the formula is meant to suggest (eb) that the probability of life beginning on a given planet is very likely connected to (e&eb) whether there are building blocks available on a given planet, and how much of them there might be.

(911) As more is learned by space researchers it is hoped that the formula could help narrow down (e) search targets by offering (eb) a statistical probability of success for (e) a given planet.

(912) The researchers note that the building blocks available do not necessarily have to be (=) of the type that led to (eb) life beginning on our planet.

**(913)** They also suggest that implied in the formula is (=) the likelihood that life beginning events are (=) more likely to occur in solar systems where (e) there are multiple planets, which allow (eb) for an opportunity to share materials that could lead to (eb)building blocks **More information: Quantifying the origins of life on a planetary scale, Proceedings of the National Academy of Sciences, www.pnas.org/cgi/doi/10.1073/pnas.1523233113**

**Abstract**

(914) A simple, heuristic formula with (e&eb) parallels to the Drake Equation is introduced to help focus discussion on open questions for (e) the origins of life in a planetary context.

(915) This approach indicates (eb) a number of areas where (e) quantitative progress can be made on parameter estimation for (e) determining origins of life probabilities, based on (e) constraints from Bayesian approaches.

(916) Authors discuss a variety of "microscale" factors and their role in determining (eb) "macroscale" abiogenesis probabilities on (e&eb) suitable planets.

(917) Authors also propose that impact ejecta exchange between (e&eb) planets with parallel chemistries and (e&eb) chemical evolution could in principle amplify (eb) the development of (e) molecular complexity and (e&eb) abiogenesis probabilities.

(918) This amplification could be very significant, and both bias conclusions about abiogenesis probabilities based on (e) the Earth and provide (eb) a major source of variance in the probability of life arising in (eb) planetary systems.

**(919)** Authors use heuristic formula to suggest (eb) a number of observational routes for improving constraints on (e) origins of life probabilities. **Read more at: http://phys.org/news/2016-07-equation-quantify-life-planets.html#jCp**

**Brian Koberlein Shared publicly - Yesterday 8:00 PM the Illusion of Truth**

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| (920) | Physical theories are often presented as (=) a description of what's really going on. |
| (921) | Forces act on (e&eb) a baseball, causing (eb) baseball to fall. |
| (922) | Atoms collide and fuse in (eb) the heart of a star, releasing heat and energy. Science is true, as is often said, and our scientific theories encapsulate this truth. But this isn't entirely true. |
| (923) | For example, quantum mechanics is a strange, sometimes confusing theory. Objects can be (=) particles and waves. |
| (924) | They do strange things like tunnel through (e&eb) barriers, and can even appear and disappear (e&eb). |
| (925) | Given all this strangeness, what's really going on? What's the true nature of quantum reality? It depends on (=) which approach you want to take. |
| (926) | In one view quantum objects are described by (e) a wave-like probability function. |
| (927) | When you interact with (e&eb) or "observe" (e) the object this "wavefunction" collapses into (eb) a definite state. |
| (928) | That's view commonly presented, but a quantum system can also be described by (e) its transitions between energy states. |
| (929) | Since this uses a mathematical method involving (e&eb) matrices, it's known as (=) matrix mechanics. |
| (930) | Both wavefunctions and matrix mechanics give (eb) the same results, but their view of what really goes on is very different. |
| (931) | Then there's the path integral method |
| (932) | . Rather than a wavefunction or matrix transitions, path integrals imagine quantum objects can take almost any path between (e&eb) two states. |
| (933) | By summing all the possible paths you can derive (eb) the odds that it will occur. |
| (934) | So which one is true? Are quantum objects distributed waves of (e) probability? |
| (935) | Are they simply transitions between (e&eb) energy states? |
| (936) | Do they take an infinite number of (e) paths between states? |
| (937) | Each of these models makes (eb) the same predictions, so one version is no more "true" than the others. |
| (938) | What happens in practice is that we'll use whatever method is useful at the time. They are equivalent models, so the best model for the job is the one we'll use. The only reason the wavefunction view is so common is that it's (=) the version usually taught to introductory students. You might think this uncertainty of truth is due to (e) the behavior of quantum physics itself. It's so strange and counterintuitive that we can't (e) wrap our puny brains around what's really going on. But the same thing occurs in lots of other fields. Even something as straight forward as basic Newtonian physics |
| (939) | Toss a baseball in the air and (eb) Earth's gravitational force pulls it down. |
| (940) | The force of gravity is a simple truth, right? While we often describe classical motion in terms of (e&eb) forces and acceleration, we can also describe it in terms of (e&eb) energy and momentum. |
| (941) | In the Lagrangian and Hamiltonian approach, the path of a baseball is (=) the optimized path among possibilities. |
| (942) | In this view a baseball's path is (=) the extrema of an energy equation, and force can be derived as a necessary consequence of (e) this. |
| (943) | In the relativistic view the baseball follows (e&eb) a geodesic, which is (=) the minimal path through (e&eb) space and time. |

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| (944) So is a baseball's motion due to (e) a gravitational force, an energy extrema, or a spacetime geodesic? |
| (945) As with quantum theory, different approaches yield (eb) the same result. They are mathematically equivalent, so we can use whatever method is most useful at the time. |
| (946) At its core, science is less about truth and more about models. The metaphysics underlying a model is useful only as far as it allows us (eb) to make better predictions, generate (eb) new ideas, or bring models together as (=) a cohesive whole |
| (947) This is why we have no problem using (e) classical gravity to calculate the path of (e) a spacecraft through (e&eb) the solar system, while (e) using special relativity to account for (e) the Doppler shift of the spacecraft's radio signals. |
| (948) It's why we can use (e) quantum physics to study atoms in the morning, and general relativity to (e) study black holes in the afternoon. |
| (949) In regimes where models conflict with each other it isn't a failure of truth, but instead shows an opportunity to develop a better model. |
| (950) It could be that with each better model we move closer to (e&eb) the truth about reality. |
| (951) The truth is out there, and science strives to move towards (e&eb) that truth. |
| (952) It's a common view, and certainly the search for truth has driven (eb) many scientists to develop better and better models. |
| (953) But the real power of science is the recognition that what we have are models. Our models can be powerful, but they are always a bit tentative. There's always a chance that they might just be an illusion of truth. |

**Synopsis: A Single-Photon Cheshire Cat July 1, 2016 Researchers detected the polarization of a photon separate from the photon itself, just as the grin of Lewis Carroll's Cheshire cat can appear apart from the cat's body Synopsis figure APS/Carin Cain**

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| (954) In a quantum Cheshire cat, one property of a particle is present detached from (e) the particle itself, just as (=) the grin of the feline in Alice in Wonderland can exist in the absence of (e) the whole animal. |
| (955) Now a research team has successfully separated (e&eb) a single photon from (e) its polarization. |
| (956) The effect was previously demonstrated (eb) with neutrons and also with light in a classical scenario, but the latest experiment uses single photons. |
| (957) To produce the Cheshire cat, **Maximilian Schlosshauer and his colleagues** at the University of Portland in Oregon used an interferometer where (e) light hits a beam splitter, goes in two, opposite directions around a circuit (but along slightly offset routes), and then recombines at (eb) the entrance. |
| (958) The incoming polarization of light was horizontal, but this was switched to (e&eb) vertical for the counterclockwise loop. |
| (959) Each path included a device to adjust (e&eb) the polarization further and another to (e&eb) scatter away a small number of photons of (e) a designated polarization, the equivalent of detecting their presence. |
| (960) The team fired single photons into (e&eb) the interferometer so that only one photon was present at (eb) a time, traversing both paths simultaneously. |
| (961) They placed a horizontal polarizer in front of the output photon detector and gradually varied (e&eb) the difference between the two path lengths. |

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| (962) This variation generated (eb) a sine-wave interference pattern from (e) the exiting photons only when the polarization was tweaked in (eb) the counterclockwise path. |
| (963) At the same time, the average number of photons counted decreased only when (e) light was scattered from the clockwise path. |
| (964) Simultaneously demonstrating these different effects for two routes of a single photon constitutes (eb) evidence of the Cheshire cat. |
| **(965)** The effect might someday improve precision experiments if (e) an unwanted property could be moved away from the property being measured. **This research is published in Physical Review A. – David Ehrenstein David Ehrenstein is the Focus Editor for Physics.** |

**Observation of the quantum paradox of separation of a single photon from one of its properties James M. Ashby, Peter D. Schwarz, and Maximilian Schlosshauer Phys. Rev. A 94, 012102 – Published 1 July 2016**

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| (966) Authors report an experimental realization of the quantum paradox of (e) the separation of a single photon from (e) one of its properties (the so-called "quantum Cheshire cat"). |
| (967) Authors use a modified Sagnac interferometer with (e&eb) displaced paths to produce (eb) appropriately pre- and postselected states of (e) heralded single photons. |
| (968) Weak measurements of photon presence and (e&eb) circular polarization are performed in each arm of (e) the interferometer by introducing (eb) weak absorbers and small polarization rotations and analyzing changes in (eb) the postselected signal. |
| (969) The absorber is found to have (e) an appreciable effect only in one arm of the interferometer, while (e) the polarization rotation significantly affects (e&eb) the signal only when (e) performed in the other arm. |
| (970) Authors carry out both sequential and simultaneous weak measurements and find (eb) good agreement between measured and predicted weak values. |
| (971) In the language of Aharonov et al. and in the sense of the ensemble averages described by (e) weak values, the experiment establishes (eb) the separation of a particle from one its properties during the passage through (e&eb) the interferometer. |

**Viewpoint: Relativity Gets Thorough Vetting from LIGO Frans Pretorius, Department of Physics, Princeton University, Princeton, NJ 08544, USA May 31, 2016• Physics 9, 52 The gravitational-wave signal observed by the LIGO detectors shows no deviation from what general relativity predicts. Figure caption expand figure Frans Pretorius, APS/Carin Cain**

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| (972) Figure 1: The signal from one of the LIGO detectors in Hanford, Washington, is shown with two representations of (e) the best-fit numerical relativity (NR) waveform. |
| (973) The filtered NR waveform illustrates how the raw waveform is perceived by (e) the detector, |
| (974) On September 14, 2015, the LIGO interferometers made the first direct measurement of (e) gravitational waves [1]. |
| (975) The historic discovery came nearly 100 years after Einstein showed that his general relativity theory predicts (eb) the existence of gravitational waves and over 50 years since (e) the first quest to observe the Universe in gravitational waves [2]. |
| (976) The observed event, labeled GW150914, was (=) the collision of two black holes, one weighing 29 solar masses, the other 36 solar masses. |

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| (977) This binary coalesced into (e&eb) a remnant black hole weighing 62 solar masses and rotating (e&eb) at 67% the maximum rate allowed (eb) for a black hole. |
| (978) The most statistically significant portion of (e) the signal came from (e) the last several gravitational-wave cycles of the collision, where (e) the most nonlinear, strong-field dynamics of the theory manifest. |
| (979) One couldn't have asked for a better first event to learn about (e&eb) this regime of general relativity. |
| (980) The first thing that GW150914 can teach us is how closely nature follows (eb) Einstein's prescriptions for (e&eb) colliding black holes. |
| (981) Researchers from the LIGO team and the Virgo team have performed a series of tests to see if (e) any violations of (e) general relativity might be lurking in the data of GW150914 [3]. |
| (982) They first determined (eb) the theoretical waveform, or "template," from (e) general relativity simulations that best fit the measured signal (see Fig. 1). |
| (983) Their main finding is that the residual signal—obtained by (e) subtracting this template from (e) the data—is consistent with (e&eb) noise. |
| (984) This is, in a sense, the "everything" test, as (e) the template folds in (e&eb) all that general relativity predicts about the merger: the decay of the last few orbits of (e) the two black holes, their violent collision (e&eb) and astonishingly rapid relaxation to (e) a quiescent rotating black hole, the propagation of (e) gravitational waves through (e&eb) the cosmos, and finally the transformation of (e) these waves into(e&eb) a chirp in the LIGO detectors. |
| (985) Given the large signal-to-noise ratio of 24, the close agreement between the template and (=) the data puts limits on (e) alternative models. |
| (986) Specifically, the predictions of these models can only disagree with (e&eb) the general relativity prediction for GW150914 by, at most, 4%. |
| (987) In addition to this comprehensive test, the LIGO/Virgo teams also performed other checks that pick apart the signal, in order to address particular aspects of general relativity and black hole physics. In one such test, they compared the early and (e&eb) late parts of the waveform. |
| (988) From the waves generated during (e&eb) the initial inspiral, one can infer (eb) the masses and spins of the binary. |
| (989) Given these numbers, general relativity allows (eb) the mass and spin of (e) the remnant black hole to be calculated. |
| (990) The researchers compared this result to (e&eb) an independent measurement of (e) the mass and spin obtained from (e) the late-time ringdown, and to within (eb) uncertainties, the two estimates were (=) consistent. |
| (991) Another test isolated coefficients in the so-called Post Newtonian (PN) series describing (eb) the inspiral. |
| (992) The PN series is (=) an expansion of the gravitational-wave signal in powers of (e) the velocity vv of the sources relative to (e&eb) the speed of light cc. |
| (993) Observations of binary pulsars, which show (eb) a decrease in their orbital period due to (e) gravitational-wave emission, have constrained (e) PN terms, especially at low order. |
| (994) For higher-order terms associated with (e&eb) strong-field emission, the LIGO/Virgo teams were able to improve (eb) the limits by many orders of magnitude compared to (e&eb) the binary pulsar constraints. |

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| (995) At first glance it might seem these new results simply add to (e) the litany of previous tests that general relativity passed with flying colors. |
| (996) However, GW150914 is vastly more significant than that. Other tests have shown that general relativity is valid in (eb)the weak-field regime, but GW150914 gives (eb) the first direct observational evidence that general relativity is the correct theory of nature governing (e&eb) the extreme spacetime dynamics around the event horizon of (e) a black hole. |
| (997) One could argue that binary pulsar observations already told us that gravitational waves exist [4], but GW150914 is more luminous in gravitational waves than pulsar binaries by (e) at least 24 orders of magnitude. |
| (998) It's true that the evolution of the Universe as a whole is driven by (e) strong gravity, but it would be difficult to argue that precision cosmological measurements can constitute (eb) a test of general relativity when (e) one of the dominant inputs, dark energy, is still a mystery. |
| (999) As far as black holes are concerned, astronomers have (e) observed ultra compact dark objects that are consistent with (e&eb) black holes. |
| (1000) The most striking data comes from (e) stellar orbits about the 4 million solar mass objects in (eb) the center of our Galaxy [5], but present observations cannot constrain it to be (=) a black hole as described by general relativity. |
| (1001) GW150914 has given us the first morsels of quantifiable evidence that the rich physics of black holes is as predicted by general relativity. For example, the LIGO observations are consistent with (e&eb) so-called cosmic censorship [6], which says that gravitational singularities must remain hidden behind an (eb) event horizon. |
| (1002) If a "naked" singularity were produced in (eb) the energetic collision of GW150914, the data would likely have shown (eb) some sort of anomaly, but in exactly what form, we don't know. |
| (1003) Another example involves the "no-hair" theorem [7], which says that the only properties of a black hole are (=) its mass, spin, and charge (though astrophysical black holes are expected to have negligible charge). |
| (1004) By consequence, these properties should determine (eb) the frequency and decay constants of (e) all the infinitely many ringdown modes of (e) a black hole. |
| (1005) With GW150914, only the least-damped ringdown mode is (=) visible above the noise, but when combined with (e&eb) the remnant mass and spin inferred from (e) the inspiral, this is enough to count as (=) the first piece of evidence for a black hole's baldness. |
| (1006) GW150914 should also be a test for (e) alternative theories of gravity or exotic alternatives to (e) black holes within general relativity, such as boson stars [8]. |
| (1007) Some of these alternative models are concerned with (e&eb) explaining dark energy or dark matter and do not (e) therefore, predict any deviation from relativity at (eb) the scale of GW150914. |
| (1008) However, other models do in fact deviate from (e) conventional relativity in their predictions of (e) the gravitational-wave signal from (e) a binary merger. |
| (1009) The largest deviations should presumably occur during (e&eb) the last several gravitational-wave cycles, where (e) the signal is strong, and deviations should be (=) easy to detect or constrain. |
| (1010) Embarrassingly though, the majority of alternatives that claim deviations can't (e) really say much about this event, as the two-body problem hasn't been solved in these theories. |
| (1011) These models can only reproduce (eb) a fraction of the full gravitational-wave signal, (eb) so the strong constraints one might have hoped for (e) cannot be leveled. |

(1012) Exotic alternatives within (eb) general relativity are in a similar boat, or worse for (e) objects like gravastars, which have (e) no well-defined theoretical framework to say what happens when (e) two such objects collide.

(1013) Regardless, what is clear from the data is that any non-black-hole solution must imagine a form of (e) exotic matter with a viscosity larger than (e) any known material to explain (eb) the rapid damping of oscillations observed during the ringdown.

(1014) The inferred event rate for heavy black hole mergers like GW150194 is high, and LIGO has another factor of 3 improvement in sensitivity slated before it reaches (eb) design sensitivity.

(1015) Within 5 years then, we can expect dozens of detected events, with the loudest having a signal-to-noise ratio upwards of (e) 100.

(1016) This will allow (eb) for exquisite insight into the nature of extreme gravity beyond (e) what GW150194 has already gifted us with. **This research is published in Physical Review Letters**.

**Tests of General Relativity with GW150914 B. P. Abbott et al. (LIGO Scientific and Virgo Collaborations) Phys. Rev. Lett 116, 221101 – Published 31 May 2016**

(1017) The LIGO detection of GW150914 provides (eb) an unprecedented opportunity to study the two-body motion of (e) a compact-object binary in (eb) the large-velocity, highly nonlinear regime, and to witness the final merger of (e) the binary and the excitation of uniquely relativistic modes of (e) the gravitational field.

(1018) Authors carry out several investigations to determine whether GW150914 is consistent with (e&eb) a binary black-hole merger in general relativity.

(1019) They find that the final remnant's mass and spin, as determined from (e) the low-frequency (inspiral) and (e&eb) high-frequency (postinspiral) phases of the signal, are mutually consistent with (e&eb) the binary black-hole solution in general relativity.

(1020) Furthermore, the data following the peak of GW150914 are consistent with (e&eb) the least-damped quasinormal mode inferred from (e) the mass and spin of the remnant black hole.

(1021) By using waveform models that allow (eb) for parametrized general-relativity violations during (e&eb) the inspiral and merger phases, authors perform quantitative tests on (e&eb) the gravitational-wave phase in the dynamical regime and they determine the first empirical bounds on (e&eb) several high-order post-Newtonian coefficients.

(1022) Authors constrain (e) the graviton Compton wavelength, assuming (eb) that gravitons are dispersed in (eb) vacuum in the same way as particles with mass, obtaining a 90%-confidence lower bound of 1013 km.

(1023) In conclusion, within our statistical uncertainties, we find no evidence for violations of (e) general relativity in the genuinely strong-field regime of gravity. **Received 26 March 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.221101 © 2016 American Physical Society**

**Focus: Background Noise of Gravitational Waves March 31, 2016• Physics 9, 33**

(1024) Assuming LIGO's recently detected black hole merger is not unusual; researchers revised upward their estimate of (e) the strength of the background noise coming from (e) distant mergers across the Universe. Figure caption expand figure Caltech/MIT/LIGO Lab (deleted due to spatial constraints)

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| (1025) Catching some waves The LIGO-Livingston interferometer consists of (e) two 4-km-long arms where (e) laser light bounces back and forth. Interference between the lasers from the two arms can signal tiny changes in (e&eb) the lengths of the arms resulting from gravi... |
| (1026) The recent detection of gravitational waves emitted from (e) a pair of merging black holes implies (eb) the existence of a large population of similar pairs, which should be generating (eb) a background of waves too faint to be detected with (e&eb) present technology. |
| (1027) Using this information, researchers now estimate that this background could be (=) roughly 10 times stronger than previously expected and might be found by (e) sensitive instruments in only a few years. |
| (1028) The background signal could help astrophysicists better understand how black hole binaries form and also pave the way to detecting the waves emitted by (e) the expanding Universe shortly after the big bang. |
| (1029) Gravitational waves were detected for the first time on 14 September 2015 by the Laser Interferometer Gravitational-Wave Observatory (LIGO). |
| (1030) By reflecting laser light between mirrors set 4 km apart, the team was able to detect (e&eb) tiny fluctuations in the distance between the mirrors as (=) the waves passed. |
| (1031) The waves were triggered by (e) an event, called GW150914; involving (e&eb) the merger of a pair of black holes having (e) masses of about 29 and 36 times that of the Sun. |
| (1032) Researchers expect that similar pairs of black holes, as well as neutron stars, should be merging constantly, creating (eb) gravitational waves that should bombard (e&eb) the Earth from all directions and make the equivalent of (=) a faint popping noise in a detector. |
| (1033) Previous estimates of the likely strength of this background had to make assumptions about (e&eb) how many pairs there might be and about the distribution of their masses. |
| (1034) These estimates remained speculative, as no actual binaries were known. |
| (1035) “What the recent observation has changed,” says Nelson Christensen of Carleton College in Minnesota and the LIGO Scientific Collaboration (LSC), “is that we now know there really is (=) a population of black hole binaries. We have a small bit of information about it.” |
| (1036) The new analysis—carried out by LSC researchers working with (e&eb) the European Virgo collaboration—estimates (eb) the energy density in the gravitational wave background in (eb) the frequency range below 100 Hz, where (e) the LIGO detector is most sensitive. |
| (1037) To make this estimate, they assumed (eb) a population of black hole binaries of various sizes, having (e) a range of distances from Earth, and created through (e&eb) a variety of scenarios. |
| (1038) They then calculated the amount of gravitational radiation that would result from (e) the merging of these objects. |
| (1039) The team then adjusted this estimate by using (e) details of the GW150914 event—especially the masses of the black holes involved—to make (eb) a most likely guess at the actual distribution of black hole binaries. |
| (1040) The idea is that the observed merger is (=) most likely representative of (e) the general population, rather than being extremely unusual. |
| (1041) If the black hole binary population is represented as (=) a bell curve, or Gaussian, the event should lie within (eb) the main bulge, not in the tails. |
| (1042) The observed mass of GW150914 suggests (eb) that the main population of black hole binaries could be (=) roughly 20 times more massive than expected, which would generate (eb) a background roughly 10 times “louder” than (e) previous estimates. |

(1043) However, Christensen emphasizes that the result has (e) lots of uncertainty, as it's based on (e) only one event.

(1044) "This estimate seems very plausible to me," says astrophysicist Maura McLaughlin of the University of West Virginia in Morgantown. She is a member of the NANOGRAV collaboration, which tries to detect gravitational waves at much lower frequencies than LIGO can observe. "I can't see any holes in their arguments. Unless LIGO got really lucky, then there must be many more of these systems than previously thought." This level of gravitational background implies (eb) that the LIGO and Virgo detectors might be sensitive enough to see it in 2–5 years, sooner than physicists previously expected, says Christensen, which has led to much anticipation. "Everyone is super excited," he says.

(1045) Detecting the background would provide (eb) a confirmation of the presence of gravitational waves carrying (e&eb) energy throughout the Universe. It would also be useful, the researchers suggest, in gaining (e) further understanding of precisely how black holes form.

(1046) And subtracting this background might someday allow (eb) detection of the fainter waves coming from (e) the expansion of the Universe moments after the big bang. This research is published in Physical Review Letters. –Mark Buchanan Mark Buchanan is a freelance science writer in Normandy, France.

**GW150914: Implications for the Stochastic Gravitational-Wave Background from Binary Black Holes B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 131102 – Published 31 March 2016**

(1047) The LIGO detection of the gravitational wave transient GW150914, from (e) the inspiral and merger of (e) two black holes with masses $\gtrsim 30 M_{\odot}$, suggests (eb) a population of binary black holes with (e&eb) relatively high mass.

(1048) This observation implies (eb) that the stochastic gravitational-wave background from binary black holes, created from (e) the incoherent superposition of (e) all the merging binaries in the Universe, could be (=) higher than previously expected.

(1049) Using the properties of GW150914, authors estimate the energy density of (e) such a background from (e) binary black holes.

(1050) In the most sensitive part of the Advanced LIGO and Advanced Virgo band for (e) stochastic backgrounds (near 25 Hz), authors predict $\Omega_{GW}(f=25\ \mathrm{Hz}) = 1.1^{+2.7}_{-0.9} \times 10^{-9}$ with 90% confidence.

(1051) This prediction is robustly demonstrated for (e) a variety of formation scenarios with (e&eb) different parameters.

(1052) The differences between models are (=) small compared to the statistical uncertainty arising from (e) the currently poorly constrained (e) local coalescence rate.

(1053) Authors conclude that this background is potentially measurable by (e) the Advanced LIGO and Advanced Virgo detectors operating at their projected final sensitivity. DOI:http://dx.doi.org/10.1103/PhysRevLett.116.131102 © 2016 American Physical Society

**Study opens new realms of light-matter interaction July 15, 2016 by David L. Chandler. Some sentences might have been omitted due to spatial constraints. Please bear with me.**

(1054) Emission spectra are a widely used method for (e) identifying chemical compounds; the bright lines reveal (eb) the different frequencies of light that can be emitted by (e) an atom.

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| (1055) Here, a normal emission spectrum for (e) an atom in a high-energy state (top) is compared to (e&eb) the emission from the same atom placed just a few nanometers (billionths of a meter) away from (e) graphene that has been doped with (e&eb) charge carriers (bottom). |
| (1056) For each energy-level transition, an orange line (or purple cloud) appears if (e) that transition is estimated to be faster than (e) one per microsecond — making it frequent enough to be (=) observed. **Credit: Massachusetts Institute of Technology** |
| (1057) A new MIT study could open up new areas of technology based on (e) types of light emission that had been thought to be (=) "forbidden," or at least so unlikely as to be (=) practically unattainable. |
| (1058) The new approach, the researchers say, could cause (eb) certain kinds of interactions between light and (e&eb) matter, which would normally take billions of years to happen, to take place instead (e&eb) within billionths of a second, under (e&eb) certain special conditions. The findings, based on a theoretical analysis, are reported today in the journal Science in a paper by MIT doctoral student Nicholas Rivera, Department of Physics Professor Marin Soljačić, Francis Wright Davis Professor of Physics John Joannopoulos, and postdocs Ido Kaminer and Bo Zhen. |
| (1059) Interactions between light and matter, described by (e) the laws of quantum electrodynamics, are the basis of a (e) wide range of technologies, including (e) lasers, LEDs, and atomic clocks. |
| (1060) But from a theoretical standpoint, "Most light-matter interaction processes are 'forbidden' by (e) electronic selection rules, which limits (e) the number of transitions between energy levels we have access to," Soljačić explains. |
| (1061) For example, spectrograms, which are used to (e) analyze the elemental composition of materials, show (eb) a few bright lines against (e&eb) a mostly dark background. |
| (1062) The bright lines represent (eb) the specific "allowed" energy level transitions in (eb) the atoms of that element that can be accompanied by (e) the release of a photon (a particle of light). |
| (1063) In the dark regions, which make up most of the spectrum, emission at those energy levels is (=) "forbidden." |
| (1064) With this new study, Kaminer says, "we demonstrate theoretically that these constraints can be lifted" using (e) confined waves within (e&eb) atomically thin, 2-D materials. |
| (1065) "We show that some of the transitions which normally take the age of the universe to happen could be made to happen within (eb) nanoseconds. |
| (1066) Because of this, many of the dark regions of a spectrogram become (=) bright once an atom is (=) placed near a 2-D material." |
| (1067) Electrons in an atom have (e) discrete energy levels, and when they hop from one level to another they give off (eb) a photon of light, a process called spontaneous emission. |
| (1068) But the atom itself is much smaller than the wavelength of the light that gets emitted—about 1/1,000 to 1/10,000 as big—substantially impairing (e&eb) the interactions between the two. |
| (1069) The trick is, in effect, to "shrink" (e) the light so it better matches the scale of the atom, as (=) the researchers show in their study. |
| (1070) The key to enabling a whole range of interactions, specifically transitions in (eb) atomic states that relate to absorbing or emitting light, is (=) the use of a two-dimensional material called graphene, in which light can interact with (e&eb) matter in the form of plasmons, a type of electromagnetic oscillation in the material. |
| (1071) These plasmons, which resemble photons but have wavelengths hundreds of times shorter, are (=) very narrowly confined in the graphene, in a way that makes some kinds of interactions with (e&eb) that matter many orders of magnitude more likely than (e) they would be in ordinary materials. |

(1072) This enables a variety of phenomena normally considered unattainable, such as (e) the simultaneous emission of multiple plasmons, or two-step light-emitting transitions between (e&eb) energy levels, the team says.

(1073) This method can enable the simultaneous emission of (e) two photons that are "entangled," meaning (eb) they share the same quantum state even when (e) separated.

(1074) Such generation of entangled photons is (=) an important element in quantum devices, such as those that might be used for cryptography.

(1075) Making use of these forbidden transitions could open up (eb) the ability to tailor the optical properties of materials in ways that had not been thought possible, Rivera says.

(1076) "By altering these rules" about the relationship between light and (e&eb) matter, "it can open (eb) new doors to reshaping (e&eb) the optical properties of materials."

(1077) Kaminer predicts that this work "will serve as (=) a founding piece for (e) the next generation of studies on light-matter interactions" and could lead to (eb) "further theoretical and experimental advances in many fields which rely on (e) light-matter interactions, including (e) atomic, molecular and optical physics, photonics, chemistry, optoelectronics, and many others."

(1078) Beyond its scientific implications, he says, "this study has possible applications across multiple disciplines, since (e) in principle it has potential to enable the full use of (e) the periodic table for optical applications."

(1079) This could potentially lead to (eb) applications in spectroscopy and sensing devices, ultrathin solar cells, and new kinds of materials to absorb solar energy, organic LEDs with higher efficiencies, and photon sources for (e) possible quantum computing devices."From the standpoint of fundamental science, this work lays the groundwork for a subfield that just a few years ago was difficult to imagine and until now was largely unexplored," Soljačić says.

(1080) "Two-dimensional materials confine (e) fields to a surface and motion to a plane, making possible (eb) many effects that are orders-of-magnitude too weak to (e) appear in a bulk volume," says Jason Fleischer, an associate professor of electrical engineering at Princeton University, who was not involved in this research.

(1081) This work, he says, "systematically explores how 2-D materials improve light-matter interactions, laying a theoretical foundation for faster electronic transitions, enhanced sensing, and better emission, including the compact generation of broadband and quantum light." More information: N. Rivera et al: Shrinking light to allow (eb) forbidden transitions on the atomic scale, Science (2016). DOI: 10.1126/science.aaf6308 Journal reference: Science search and more info website: Provided by: Massachusetts Institute of Technology search and more info website. This story is republished courtesy of MIT News (web.mit.edu/news office/), a popular site that covers news about MIT research, innovation and teaching.

**Viewpoint: Peering Back to the Beginning of Time Lawrence M. Krauss, Department of Physics and School of Earth and Space Exploration, Arizona State University, Tempe, AZ 85287, USA and Mount Stromlo Observatory, Research School of Astronomy and Astrophysics, Australian National University, Weston, ACT 2611, Australia June 19, 2014• Physics 7, 64**

(1082) The BICEP2 collaboration reports the detection of BB-mode polarization of (e) the cosmic microwave background—a signal that might originate from (e) gravitational waves created by (e) inflation during the very earliest moments in the evolution of our Universe. Figure caption expand figure P. A. R. Ade et al. [1]; Image on homepage: BICEP2 Collaboration (deleted due to restrictions on space)

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| (1083) Figure 1: (Top) BB-mode polarization pattern observed by (e) the BICEP2 telescope [1]. |
| (1084) The twisting pattern in the polarization of (e) the cosmic microwave background (CMB) might originate from (e) gravitational waves created by inflation in the very first moments of the universe. |
| (1085) Inflation—the hypothesis that the Universe underwent (e&eb) a phase of superluminal expansion in a brief period following the big bang—has (e) the potential of explaining, from first principles, why the Universe has (e) the structure we see today. |
| (1086) It could also solve outstanding puzzles of standard big-bang cosmology, such as why the Universe is, to a very good approximation, (=) flat and isotropic (i.e., it looks the same in all directions). |
| (1087) Yet we do not yet have a compelling model, based on (e) fundamental particle physics principles, that explains inflation. |
| (1088) And despite its explanatory power and a great deal of suggestive evidence, we still lack an unambiguous and direct probe of (e) inflation. |
| (1089) Theorists have developed different models for inflation, which all share (e&eb) a common, robust prediction: Inflation would have created (eb) a background of gravitational waves that could have (eb) an observable effect. |
| (1090) These waves would cause (eb) subtle, characteristic distortions of (e) the cosmic microwave background (CMB)—the oldest light in the Universe, released when (e) photons decoupled from matter and the Universe became (=) transparent to radiation |
| (1091) Until now, such a background of gravitational waves has (e) escaped detection. It is therefore very exciting that the BICEP2 (Background Imaging of Cosmic Extragalactic Polarization) collaboration has now reported in Physical Review Letters the observation of polarization modes of CMB radiation [1] that are consistent with (e&eb) such a background. |
| (1092) If confirmed, BICEP2's finding would provide (eb) the first direct probe of physical processes that occurred a mere 10−3610-36 seconds after (e) the big bang. |
| (1093) The result also offers a rare opportunity to directly test theoretical ideas, including inflation and grand unification (the theory that merges the electromagnetic, weak, and strong forces into a single force), at energies 13 orders of magnitude greater than (e) we can currently access at the Large Hadron Collider. |
| (1094) Inflation is based on (e) the possibility that a scalar field in the early Universe got momentarily stuck in (eb) a "false vacuum state"—a metastable condition in which the Universe wasn't at (eb) its global potential minimum. |
| (1095) The energy stored in this vacuum state could drive (eb) an exponential expansion of spacetime (similar to, but much faster than, the expansion due to dark energy seen today) before (e) a phase transition to the true vacuum state occurs, releasing (eb) this energy into (e&eb) particles and radiation. |
| (1096) This creates (eb) an initially hot state for (e) a subsequent big-bang-like expansion from (e) which our Universe subsequently evolved [2]. |
| (1097) If the Universe grew by at least a factor of 10261026, (eb) inflation can explain (eb) both the apparent flatness and isotropy of the Universe [2]: |
| (1098) The expansion would stretch (eb+ eb-dot) any initial curvature of (e) the Universe to near flatness, and smooth out (e&eb) any initial inhomogeneities, just as stretching (eb+ ebdot) a piece of fabric would do. |
| (1099) Isotropy is justified by the fact that the Universe, before inflation, was once much smaller and could have been contained in (eb) a causally connected region that was in thermal equilibrium, resulting in (eb) the approximately uniform, isotropic temperature distribution observed on (eb) large scales today. |

(1100) Importantly, the theory can explain the (eb) origin of structure in the Universe: Microscopic quantum fluctuations associated with (e&eb) the scalar field driving inflation could have been "frozen in" during the exponential expansion, resulting, (eb) after inflation, in density fluctuations that might have seeded all observed large-scale structures in the Universe, including (e) all visible galaxies and clusters of galaxies.

(1101) The signs of such primordial fluctuations have indeed been detected through (e&eb) their effect on (e&eb) the CMB, first by the Cosmic Background Explorer (COBE) [3] and later by other missions.

(1102) These observational studies detected small temperature variations (anisotropies) in (eb) the CMB and found a spectrum compatible with (e&eb) that predicted by inflation.

(1103) The magnitude of these fluctuations suggests (eb) an energy scale of inflation (loosely speaking, the energy stored throughout space in (eb) the scalar field potential that drove inflation) close to (e) 10161016 giga-electron-volts (GeV).

(1104) But while consistent with inflation, such fluctuations are not (e) conclusive evidence: Their predicted shape and magnitude is known to vary with (e&eb) the details of the inflationary potential (whose form is not yet experimentally constrained).

(1105) Hence many different observations could be accommodated, and any agreement between theory and observation is at best suggestive. There is also a possibility that other, noninflationary sources might contribute. However, gravitational waves—the propagating ripples in spacetime predicted by Einstein in 1916—might provide (eb) direct evidence of inflation.

(1106) During inflation, gravitational fields would also fluctuate because of (e) quantum effects.

(1107) These fluctuations would similarly get "frozen in" and re-emerge after inflation as (=) classical gravitational waves, whose amplitude depends only on (=) the energy stored in the field driving (eb) inflation, and not on other model-dependent details [4].

(1108) If the scale of inflation were indeed as high as a few times 10161016 GeV, this background could produce (eb) specific temperature anisotropies in the CMB [5].

(1109) These anisotropies have thus far not been detected by (e) the most recent searches by the WMAP and PLANCK satellites, and their absence was used to (e) put upper limits on the scale of inflation somewhat below 10161016 GeV.

(1110) But a more sensitive CMB probe of (e) gravitational waves exists.

(1111) Large-wavelength gravitational waves produced by (e) inflation would have alternatively compressed and expanded (e&eb) space, producing (eb) spatial variations in (eb) the intensity of radiation bombarding free electrons in (eb)the early Universe (see, e.g., Ref: [6]).

(1112) As photons decoupled from (e) matter and CMB radiation was released, scattering from (e) the undulating sea of electrons would have created (eb) patterns in the polarization of CMB.

(1113) Such polarization anisotropies are predicted in (eb) inflationary models to be (=) most visible when (e) looking at the sky on spatial scales comparable to (e&eb) or somewhat larger than the size of (e) the Universe's horizon at (eb) the time when the CMB was created.

(1114) This corresponds to (e&eb) about 1 degree of the field of view of the sky.

(1115) Over the past decade, various CMB polarization experiments have thus chased the observation of polarization anisotropies on (e&eb) these scales but, so far, also yielded (eb) merely upper limits on a possible gravitational background.

(1116) All of this has now changed with (e&eb) the BICEP2 collaboration's [1] report of (e) a signal that has the predicted (eb) characteristics of such a background at precisely the angular scales where (e) the inflationary signal is expected (see Fig. 1, top deleted due to spatial restraints).

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| (1117) The BICEP2 Collaboration uses (e) a microwave telescope located at the South Pole, with (e&eb) an array of polarization-sensitive detectors pointing at (eb) a small field (a solid angle of about 11 percent of the full sky) with a resolution of approximately 0.50.5 degrees, thus sufficient to probe the relevant angular scales for the detection of (e) polarization anisotropies originating from (e) **inflation**. |
| (1118) BICEP2's search is based on (e) two expected features of primordial gravitational waves. |
| (1119) First, gravitational waves would produce (eb) a characteristic spatial polarization pattern on CMB—a consequence of (e) their quadrupole nature. |
| (1120) One can decompose (e&eb) any spatial polarization pattern into (e&eb) two orthogonal components, a "curl-free" component and (e&eb) a component involving a twisting pattern. |
| (1121) In analogy to electromagnetic modes, these polarization patterns are called "EE modes" and (e&eb) "BB modes," respectively (see Fig. 1). |
| (1122) Scalar fluctuations in the density of (e) matter would produce (eb) EE modes, while (e) gravitational waves would instead generate (eb) a BB-mode pattern. |
| (1123) The second feature is the aforementioned characteristic spatial scale of (e) expected CMB polarization signatures, including (e) wavelength scales comparable to (e&eb) or larger than the horizon scale at the time the CMB was created. |
| (1124) Since inflation would produce (eb) both gravitational waves and (e&eb) scalar density perturbations, BICEP2, like all CMB-polarization experiments, quote their results in terms of (e&eb) the ratio rr of the amplitude of (e) tensor perturbations (associated with (e&eb) gravitational waves) to (e) the amplitude associated with (e&eb) scalar density fluctuations (reflected in (eb) the magnitude of CMB temperature fluctuations measured by (e) experiments like COBE). |
| (1125) A nonzero rr implies (eb) the detection of a BB-mode pattern and primarily measures (eb) the strength of any possible gravitational wave signal which, if (e) it comes from inflation, is then determined by (e) its energy scale (e.g., see Ref. [7] for a review). |
| (1126) BICEP2 reported a nonzero value of r=0.2(+0.07−0.05)r=0.2(+0.07-0.05), on scales ranging from (e) the horizon scale up to 4–5 times larger, with (e&eb) a claimed significance of better than 5σ5σ. |
| (1127) If confirmed, a value of 0.2 corresponds to (e&eb) an energy density during (e&eb) inflation of ∼1.2×1016∼1.2×1016 GeV. |
| (1128) The claimed signal is far larger than expected, partly because (e) the PLANCK satellite, despite not having yet reported its polarization measurements, has (e) previously put an upper limit on (eb) rr (<0.11<0.11). PLANCK's estimate is based on (e) the absence of a gravitational wave signature on (eb) the CMB temperature anisotropy on (eb) its largest scales (PLANCK observes the whole sky and therefore angular scales up to 360360 degrees). |
| (1129) The apparent discrepancy between the PLANCK upper bound and (e&eb) the BICEP2 detection might be a cause for concern. |
| (1130) However, the two bounds probe different spatial scales, and the BICEP collaboration has noted that, while tensor perturbations are expected to be (=) roughly scale independent, some scale dependence of (e) scalar density perturbations could relax (e&eb) the PLANCK constraint and bring (eb) the two results into agreement. |
| (1131) Of greater concern, of course, is the possibility that the signal could be determined by (e) foreground contamination—any source along (e&eb) the line of sight of the detector that could distort or contaminate (e&eb) the primordial CMB signal. |
| (1132) BICEP2 tested for (e) a wide variety of possible foregrounds, claiming (eb) they cannot be responsible for (e) the observed signal. |

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| (1133) One obvious candidate arises from (e) gravitational lensing: The intervening mass of (e) astrophysical bodies could distort (e&eb) the spatial pattern of polarization in a way that that could transform (e&eb) an EE-mode signal into (e&eb) a BB-mode signal. |
| (1134) The BICEP collaboration has analyzed this possibility extensively. Based on (e) measurements of (e) CMB temperature (and hence density perturbations) and on (e&eb) models of the mass distribution, they argue that such effects would only be (=) significant on smaller angular scales than those contributing to their measurements of rr. |
| (1135) But since BICEP2's first announcement on March 17, 2014, a number of scientists have raised another foreground concern due to (e) the possible impact of (e&eb) our Galaxy's dust. |
| (1136) BICEP2 modeled dust contamination scenarios based on (e) a variety of theoretical models, using data from (e) its predecessor BICEP1 and preliminary dust estimates from (e) the PLANCK satellite. |
| (1137) They concluded that dust could not (e) reproduce the magnitude of the observed signal. |
| (1138) However, PLANCK has in the meanwhile released (eb) more complete measurements of variation in polarized dust foregrounds in (eb) many regions of the Galaxy. |
| (1139) Their analysis increases the uncertainty associated with (e&eb) the dust contribution to (e) the BICEP2 result and also accommodates (e) a potentially larger polarized dust contribution. |
| (1140) Some researchers argued that dust polarization at a level 2–3 times higher than that assumed by BICEP2 is possible and could therefore plausibly be responsible for (e) the entire signal (e.g., see Ref: [8]). |
| (1141) While the BICEP2 collaboration continues to maintain that this is extremely unlikely based on (e) their estimates, many other researchers are more skeptical. |
| (1142) In any case, a variety of upcoming measurements, from PLANCK and other detectors are (=) expected to resolve this important ambiguity within the coming year. |
| (1143) Even if the BICEP2 signal does arise from (e) primordial gravitational waves, various consistency checks will be required (e) before it can be definitively argued that the signal arises from (e) inflation. |
| (1144) Inflation is certainly the best-motivated possibility, and the spectral shape of the signal, at least for (e) the five largest-scale modes observed by BICEP2, agrees extremely well with (=) inflationary predictions (see Fig. 1, bottom). |
| (1145) Several results have already appeared that strongly constrain (e) other possible sources, ruling out (e) for example, the possibility of gravitational waves generated by (e) other possible phase transitions after inflation [9] |
| (1146) . This conclusion is reinforced by more careful analyses of the scale dependence of (e) the predicted signal (e.g., Ref. [10]). |
| (1147) But it will nevertheless be important to assure that all possible alternative sources of (e) gravitational waves are ruled out. |
| (1148) If it can be established that the BICEP2 represents (eb) a gravitational wave signal from (e) inflation, there is one additional exciting implication. |
| (1149) Standard calculations suggest (eb) that the gravitational wave signal from inflation is due to (e) quantum fluctuations in the gravitational field, and this appears to require (e) the existence of gravitons—the hypothetical elementary particles that mediate (e&eb) the force of gravity. |
| (1150) Research by Frank Wilczek and myself, based on (e) a dimensional analysis argument [11], suggests that,(eb) independently of the method used to calculate the gravitational wave spectrum, a quantum gravitational origin is (=) required. |
| (1151) If so, the BICEP2 result then implies (eb) that gravity is ultimately a quantum theory, a result of fundamental importance for physics. |

(1152) It is hard to imagine a more important or wider-ranging cosmological observation. If confirmed by upcoming experiments, the BICEP2 results will push the frontiers of empirical physics forward in an unprecedented way, offering (eb) the first direct observation of gravitational waves and of quantum gravity effects as well as a new probe of the earliest moments in the Universe's life and of physics at grand-unified energy scales. This will thus constrain both our understanding of fundamental interactions and the evolution of our Universe as a whole.

**Detection of B-Mode Polarization at Degree Angular Scales by BICEP2 P. A. R. Ade et al. (BICEP2 Collaboration) Phys. Rev. Lett 112, 241101 – Published 19 June 2014**

(1153) Authors report results from the BICEP2 experiment, a cosmic microwave background (CMB) polarimeter specifically designed to search for (e) the signal of inflationary gravitational waves in (eb) the B-mode power spectrum around $\ell \sim 80$.

(1154) The telescope comprised a 26 cm aperture all-cold refracting optical system equipped with (e&eb) a focal plane of 512 antenna coupled transition edge sensor 150 GHz bolometers each with temperature sensitivity of $\approx 300\ \mu K\sqrt{CMB s}$. BICEP2 observed from (e) the South Pole for three seasons from 2010 to 2012.

(1155) A low-foreground region of sky with (e&eb) an effective area of 380 square deg was observed to (e) a depth of 87 nK deg in Stokes Q and U.

(1156) In this paper authors describe the observations, data reduction, maps, simulations, and results. They find an excess of B-mode power over the base lensed-ΛCDM expectation in the range $30<\ell<150$, inconsistent with (e&eb) the null hypothesis at (eb) a significance of $>5\sigma$.

(1157) Through jackknife tests and simulations based on (e) detailed calibration measurements authors show that systematic contamination is (=) much smaller than the observed excess.

(1158) Cross correlating against WMAP 23 GHz maps we find that Galactic synchrotron makes a negligible contribution to the observed signal. We also examine a number of available models of polarized dust emission and find that at their default parameter values they predict power ~(5–10)× smaller than the observed excess signal (with no significant cross-correlation with our maps). However, these models are not sufficiently constrained by external public data to exclude the possibility of dust emission bright enough to explain the entire excess signal. Cross correlating BICEP2 against 100 GHz maps from the BICEP1 experiment, the excess signal is confirmed with $3\sigma$ significance and its spectral index is found to be consistent with that of the CMB, disfavoring dust at $1.7\sigma$. The observed B-mode power spectrum is well fit by a lensed-ΛCDM+tensor theoretical model with tensor-to-scalar ratio $r=0.20+0.07-0.05$, with $r=0$ disfavored at $7.0\sigma$. Accounting for the contribution of foreground, dust will shift this value downward by an amount which will be better constrained with upcoming data sets. Received 4 April 2014 DOI:http://dx.doi.org/10.1103/PhysRevLett.112.241101 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**(1159) Physicist claims to have observed quantum effects of Hawking radiation in the lab for the first time August 16, 2016 by Bob Yirka report**

(1160) This computer-simulated image shows a supermassive black hole at the core of a galaxy. The black region in the center represents the black hole's event horizon, where no light can escape the massive object's gravitational grip. The black …more

(1161) (Phys.org)—Jeff Steinhauer, a physicist at the Israel Institute of Technology, has published a paper in the journal Nature Physics describing experiments in which he attempted to create a virtual black hole in the lab in order to prove that Stephen Hawking's theory of radiation emanating from black holes is correct —though his experiments are based on sound, rather than light. In his paper, he claims to have observed the quantum effects of Hawking radiation in his lab as part of a virtual black hole—which, if proven to be true, will be the first time it has ever been achieved.

(1162) For many years, scientists believed that nothing could ever escape from a black hole. But in 1974, Stephen Hawking published a paper suggesting that something could—particles that are now called Hawking radiation. His idea was that if a particle (and its antimatter mate) appeared spontaneously at the edge of a black hole, one of the pair might be pulled into the black hole while the other escaped, taking some of the energy from the black hole with it—which would explain why black holes grow smaller and eventually disappear. Because such emissions are so feeble, no one has been able to measure Hawking radiation, so researchers have instead tried to build virtual black holes in labs to test the theory. One type of virtual black hole was proposed back in 1981 by Bill Unruh with the University of British Columbia—he suggested that an analogue might be created using water instead of light. He imagined a phonon existing at the edge of a waterfall—as the water speeds up, it begins to move faster than the speed of sound, causing it to be trapped. But if the phonon had an entangled mate that eluded the fall by moving away before getting caught up, it could escape. In this new effort, Steinhauer has built a device based on that idea and in so doing, claims he has observed an analogue of Hawking radiation.

(1163) The experiment consisted of creating an entangled pair of phonons sitting inside a bit of liquid that had been forced (via laser) to move very fast and then observing the action as one of the pair was pulled away as the liquid began to move faster than the speed of sound, while the other escaped—the fluid was a Bose-Einstein condensate of rubidium-87 atoms. After repeating the experiment 4,600 times Steinhauer became convinced that the particles were entangled, a necessity for a Hawking radiation analogue. His findings do not prove Hawking's theory to be true, of course, but they do appear to add a degree of credence that other researchers have thus far not been able to achieve. More information: Jeff Steinhauer. Observation of quantum Hawking radiation and its entanglement in an analogue black hole, Nature Physics (2016). DOI: 10.1038/nphys3863

**(1164) Abstract**

We observe spontaneous Hawking radiation, stimulated by quantum vacuum fluctuations, emanating from an analogue black hole in an atomic Bose–Einstein condensate. Correlations are observed between the Hawking particles outside the black hole and the partner particles inside. These correlations indicate an approximately thermal distribution of Hawking radiation. We find that the high-energy pairs are entangled, while the low-energy pairs are not, within the reasonable assumption that excitations with different frequencies are not correlated. The entanglement verifies the quantum nature of the Hawking radiation. The results are consistent with a driven oscillation experiment and a numerical simulation. Journal reference: Nature Physics search and more info website

Read more at: http://phys.org/news/2016-08-physicist-quantum-effects-hawking-lab.html#jCp

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as

the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t - T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

These equations can be written for

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(7) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.

(8) However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].

(9) Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.

Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

**TIME DELAY**

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period,

which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types system takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t - T_7)F_7(R_7(t - T_7) - N_7(t)M_7((R_7(t) \qquad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7}+\frac{N_7 dG_7}{dR_7}-\frac{dB_7}{dR_7}>\frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7}+\frac{dM_7}{dR_7}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also normalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_7$) and the other derived from the non-delayed or "immediate" processes($I_7$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_7=\frac{N_7 G_7}{2M_7}\frac{d(7nF_7)}{dR_7}(\geq 0) \quad (7a)$$

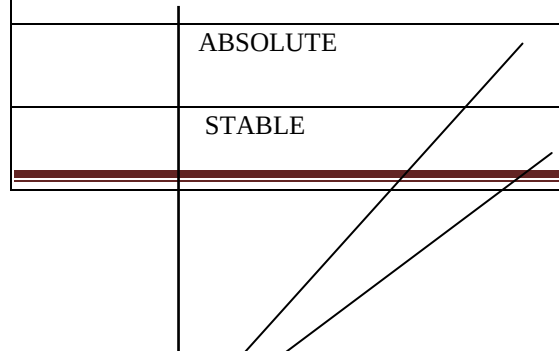
$$I_7=-\frac{N_7 G_7}{2M_7}\frac{d(7nM_7}{dR_7}(\geq 0) \quad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7>L_7-I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

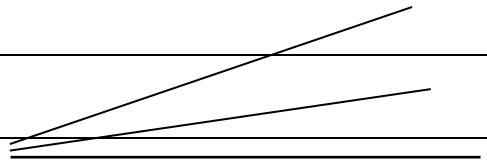
$$P_7>L_7-I_7, \quad \text{when} \quad P_7^2/4>I \quad (8a)$$

$$P_7>L_7-I_7+\frac{(I_7-P_7^2}{(L_7-I)} \quad \text{when} \quad P_7^2/4<I. \quad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,





**DELAYED PREDATOR CREDIT TYPES EFFECT WITH ITS WIDE RANGING RAMIFICATION**

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

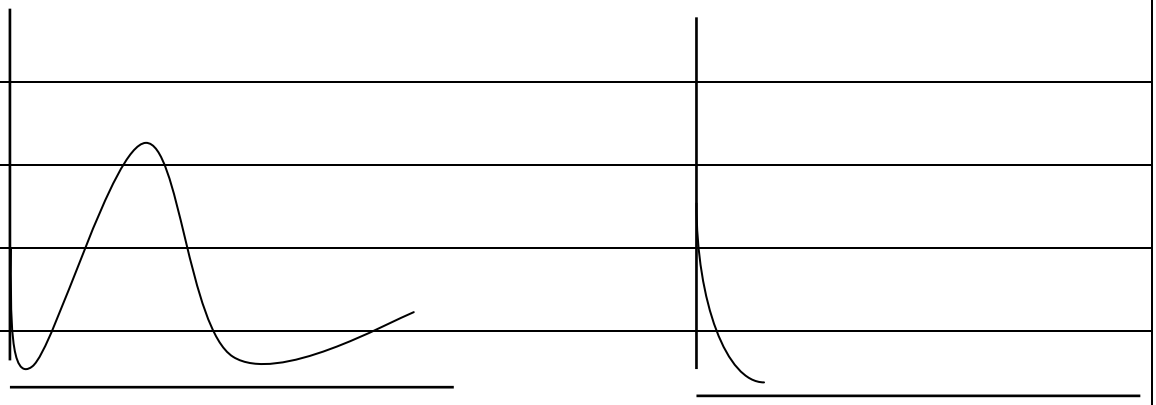
It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay

and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



**TIME DELAY**

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

To scores of problems this relationship is important. Of them some of the examples are as follows:

**New look at the galaxy-shaping power of black holes**

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| (257) | Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving. |
| (258) | "These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science. |
| (259) | "Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy." |
| (260) | This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return. |
| (261) | The remaining matter gets absorbed into the black hole, adding to (e) its mass. |
| (262) | The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne. |
| (263) | The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution. |
| (264) | Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form. |
| (265) | "Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara. |
| (266) | "It creates (eb) these bubbles, and the bubbles keep the plasma hot. |
| (267) | That's what prevents galaxies from (e) becoming even bigger than they are now." |
| (268) | Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space. |
| (269) | In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene. |
| (270) | Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve. |
| (271) | "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara. |
| (272) | "It's like nature's thermostat, which keeps these galaxies from (e) growing. |
| (273) | If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars." |

**A new look at the galaxy-shaping power of black holes**

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| (274) | An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the |
| (275) | McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said. |

(276) Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.

**Doomed satellite missions**

(277) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.

(278) At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.

(279) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.

(280) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.

(281) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.

(282) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people

(283) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."

(284) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.

(285) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(286) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.

(287) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.

(288) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.

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| (289) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses. |
| (290) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well. |
| (291) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma. |
| (292) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6. |
| (293) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus. |
| (294) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (295) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (296) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (297) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (298) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (299) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (300) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (301) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (302) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (303) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (304) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (305) We know dark matter exists only from (e) its gravitational effects. |
| (306) It holds together galaxies and (e&eb) galaxy clusters. |

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| (307) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (308) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |
| (309) They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (310) X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (311) They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (312) This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (313) An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |
| (314) Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (315) "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (316) The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (317) Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (318) Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (319) "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (320) Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (321) "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New |

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| York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (322) Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories. |
| (323) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen. |
| (324) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer. |
| (325) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants. |
| (326) These were treated as (=) "aware" and "unaware" trials, respectively. |
| (327) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it. |
| (328) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another. |
| (329) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion. |
| (330) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness. |
| (331) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said. |
| (332) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness." |
| (333) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner. |
| (334) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain. |
| (335) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said. |
| (336) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

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| (337) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter. |
| (338) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined. |
| (339) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions. |
| (340) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda. |
| (341) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general. |
| (342) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole. |
| (343) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes. |
| (344) Noda and his colleagues instead studied the change in emissions over (e&eb) time. |
| (345) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection. |
| (346) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1). |
| (347) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole. |
| (348) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved. |
| (349) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp |

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(350) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking “palm tree” organization.

(351) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society

(352) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible?

(353) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem.

(354) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other.

(355) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible.

(356) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges.

(357) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century.

(358) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching.

(359) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom.

(360) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth.

**Spheres on trees**

(361) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration.

(362) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement.

(363) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum.

(364) By visualizing the problem in this way, the researchers noticed that these particular graphs don’t (e) have lots of branches extending from lots of (e) other branches.

(365) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure.

(366) The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network.

(367) As a result, it provides (eb) an important clue for finding the global minimum.

(368) It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum.

(369) This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large.

**Thompson's problem 2**

(370) By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society

(371) The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima.

**(372)** The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node.

**Implications for Thomson's and Smale's problems**

(373) Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before.

(374) "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around."

(375) As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem.

(376) "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view."

(377) In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges.

(378) From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear.

(379) "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK.

(380) "For larger numbers, there will likely be (=) multiple funnels.

(381) Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum."

(382) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.

(383) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.

(384) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t-T_7)F_7\big(R_7(t-T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t-T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) and the effective prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this

rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Look at the resemblance of String Theory and Abhinava's earlier theory and how the model could be applied: The world of manifestation and differentiated perceptions (Yikalpa) thus extends from one Centre to (eb) the next. **What is meant again is the realisation or the evolution of individual consciousness or general ledger to be (=)the cosmic general ledger is the sole criterion for the Satchidananda state where all dualities cease** Although it is never in fact divorced from (e) the subject who resides there, the ignorant fail (e) to grasp this fact and so, cut off from (e) the Centre, the world of objectivity becomes (=) for them the sphere of Maya. Bhagavatotpala quotes the Light of Consciousness amvitprakdsa): This ever pure experience (suddhanubhava) is variegated by (e) each form [revealed within it] Even so it remains unstained (nirmala) when moving to another. Just as a cloth which is naturally white, once dyed, cannot (e) change colour without [first] becoming white again, similarly the pure power of awareness, (citi) once coloured by (e) form, is pure [again] at the Centre where that form is (=) abandoned and from whence it proceeds to (e&eb) another. In his Essence of Vibration (Spandasarpdoha), K§emaraja explains that the rise and fall of every individual perception in the field of awareness is (=) a specific pulsation of consciousness. This is exactly the individual general ledger transaction lifted to the transaction in cosmic general ledger. During The Initial Instant Of Perception, T Consciousness Is Manifestly Apparent And The Yogi, Participating In Its Plenitude, Observes The Outer World Without Being Attached To Any Particular Divine Body And Sacred Circle Of The Senses Singling It Out From Any Other, Like A Man Who Observes A City From A High Mountain Peak.

He Sees The Outer World Reflected Within His Consciousness Free Of Thought-Constructs And So 'Stamps* The Outer On The Inner While Absorbing The Object And Means Of Knowledge In The Pure Subject Which Grasps Them As The Expansion Of His Own Nature. K§emaraja Says: By Penetrating Into Bhairavimudra, The Yogi Observes The Vast Totality Of Beings Rising From, And Dissolving Into, The Sky [Of Conscious- Ness], Like A Series Of Reflections Appearing And Disappearing Inside A Mirror. Through (e&eb) The Practice Of Bhairavimudra The Yogi Realises That He Is (=) The Substratum Consciousness (Adhisfhatr) Which Both Underlies (e&eb) And Is (=) The Essence Of All Things. He Discovers That Phenomena Have (e) No Independent Existence Apart From (e) Him And So Are, In This Sense, Void. At The Same Time, He Realises That Because (e) All Things Are Consciousness, They Are (=) Far From Unreal. He Views The Outer World Yet Sees (e) It Not. Beyond Both Voidness And Non-Voidness He Penetrates Into (e&eb) The Supreme Abode (Pararn Padam) Of Siva's Consciousness. [The Powers Of The Senses] Endowed With (e&eb) The Attributes Of The Great Union [Between Subject And (e&eb) Object] Whose Form Is The Wakening Of Man's Spiritual Potential (Kuntfalinf), Fill (e&eb) [With Consciousness] The Outer Clatter Of Diversity (Bhedatfambara) Born Of (e) Its Intense Power And Are Then Established In (eb) The Unobscured Abode Of The Void Of Consciousness To Shine [There] Eternally. Thus Residing Beyond (e) Being And Non-Being, The Sole Protector Of The Unity Which Is Tranquil And Expanding [Consciousness], Whose Glory Is All-Embracing And Form Unobscured, Is Called Bhairavimudra. Through (e&eb) The Practice Of Bhairavimudra, The Yogi Unites The Universal Vibration Of T Consciousness With (e&eb) The Individual Pulsation Of Objectivised 'This' Consciousness. The Two Aspects Of Consciousness Are Now In (eb) A State Of Equilibrium Like (=) The Two Pans Of An Evenly Weighed Balance And The Yogi Experiences (e&eb) The Pure Knowledge (Suddhavidya) That: 'I Alone Am (=)All Things'. Thus Becoming (=) The Master Of The Wheel Of Energies He Is (=) Free, Like Siva, To Create And Destroy. When [The Yogi] Is Well Established, Without (e) Wavering, Solely In The Integral Egoity Of His Authentic Nature, The Spanda Principle, And Is

Absorbed In Contemplation (Samavisfa), He Becomes (=) One With It The Doctrine Of Vibration (Tanmaya). Then . . . Dissolving And Creating The Universe By Means Of (er) His Introverted And Extroverted Absorption, He Destroys And Creates (e&eb) All Things Out Of Sarikara, His Innate Nature. [Thus] He Assumes The State Of The Universal Experiencer And Having Absorbed All That Is To Be Experienced From [The Grossest Level] — Earth — To (The Subtlest) — Siva — He Reaches (eb) The State Of The Supreme Subject By (e) Progressively Recognising [His Identity With Him]. Thus, Introverted And Extroverted Absorption Both Lead To (eb) The Recognition Of The Pulsation (Spanda) Of (e) One's Own Consciousness. This Is Where (e) Individual General Ledger Becomes (=) Cosmic General Ledger. At The Level Of Consciousness Corresponding To (e&eb) Siva's Basic State (Saryibhavdvastha), The Alternation From Inner To (e&eb) Outer Is Instantaneously Resolved Into (e&eb) The Vibration Of His Nature. When The Yogi Finally Comes To Be (=) Constantly Aware Of This Reality, His Enlightenment Is (=) Full And Perfect. Freed Of All Means (Anupaya) And Delighting In The Power Of His Bliss (Anandasakti), He Knows And Does Whatever He Pleases. The Yogi Seeking Self-Realisation Must Acquire (e) Mastery Over This Movement. Ksemaraja Stresses That The Doctrine Of Vibration Teaches That Liberation Can Only Be Achieved By (e) First Withdrawing All Sense Activity In Introverted Contemplation (Nimilanasamadhi) To (e) Then Experience The 'Great Expansion' (Mahdvikdsd) Of (e) Consciousness While (e) Recognising This To Be (=) A Spontaneous Process Within (eb) It. This Is Done Through (e&eb) The Practice Of Kramamudra. A Passage From The Now Lost Kramasiitra Explains: Although The Adept's Attention [May Be] Outwardly Directed, (eb) He Enjoys Contemplative Absorption Through (e&eb) The Introverted Aspect Of (e) Kramamudra. Initially He Turns Inward From (e)The Outside World And [Then] From Within [Himself] He Exits Into The Outer World Under The Influence Of His Absorption. There Is No Necessity Of Any Glorification Or Mortification Here. It Is Just Like Joining A Job And Retiring And Realising All That Happened Is Just By His Own Effort And Nothing Else. Thus The Sequence (Krama) In This Attitude (Mudra) [Ranges Through] Both Inner And Outer. The Yogi Must Pervade

The Surface Level Of Awareness (Yyutthana) With The Same Bliss He Experiences Plunging Into The Depths Of Con- Templative Absorption (Samadhi). Note Here That Samadhi Is A Death Like State And Still Being Matter And Information Stored You Experience Consciousness. Submerging Himself And Emerging Repeatedly From Samadhi, He Eventually Recognises That The Unity Of Consciousness Pervades Both States: The Best Of Yogi's, Who Has Achieved A State Of Complete Absorption Even When Risen From Meditation, [Inwardly] Vibrating Like A Drunkard In Blissful Inebriation From The After-Effects Of The Nectar Of Con- Templation, Sees All Things Dissolving In The Sky Of Consciousness Like A Cloud In The Autumn Sky. He Plunges Repeatedly within Himself Divine Body and Sacred Circle of the Senses and Becomes (=) Aware Of His Identity With (e&eb) Consciousness By (e) The Practice Of Introverted Contemplation. Thus Even When He Is Said To Have Risen From (e) Absorption, He Is (=) One With [His] Experience Of It. Ksemaraja Goes On To Explain That This Practice Is Called 'Mudra' Because (e) It Both Fills The Adept With (e&eb) Bliss (Mud) And Is Itself The Bliss Of (e) Consciousness. Moreover, It Dissolves (e) Away (Dra) All Bondage And 'Stamps' The Universe Of Experience With (e&eb)The Seal (Mudra) Of The Fourth State (Turiya) Of Enlightened Consciousness Beyond, And Including, The Three States Of Waking, Dreaming And Deep Sleep. It Is Called 'Krama' Because (e) It Is The Root Source Of (e) All Emanation And All Other Conscious Processes Which Succeed (e&eb) One Another In Ordered Sequence (Krama) And Is, (=)At The Same Time, Their Successive (Krama) Appearance As Well. By The Practice Of Kramamudra The Opposites Fuse And Siva And Sakti Unite. They Yogi Comes To Experience The Simultaneous Pervasion Of All The Lower, Grosser Categories Of Existence By The Higher And The Presence Of The Lower In The Higher. Commencing His Practice In A Low Form Of Bhairavimudra, The Yogi Conjoins The Outer With The Inner; Then, In Kramamudra, He Fills Both The Outer With (e&eb) The Inner And (e&eb) The Inner With (e&eb)The Outer. When He Achieves (eb) Perfection In This Two-Fold Movement, He Attains To (e) The Highest Form Of Bhairavimudra In Which The Two Merge Completely In (eb)The Experience Of (e) The Absolute (Anuttara), Free

Of (e) All Differentiation And (e&eb) Polarities. If He Fails To Maintain Awareness Of This State, He Again Falls Into (eb) Kramamudra Until He Has Finally Completely Merged All The Highest States In (e&eb) The Lower And The Lower In (e&eb) The Higher. He Then No Longer Needs To (e) Resort To (e) Any Means (Anupaya) To (e) Achieve Liberation. All He Says or Does Anything He Perceives or Thinks, Instantly Occasions In (eb) Him The Highest Level Of Consciousness. Thus The Fruit Of Bhairavimudra Is (*=)The Wonder (Camatkara) Or Amazement (Vismaya) That Overcomes (e)The Yogi When (e) He Reaches The Plane Of (e) Union (Yogabhumika) Y Xsl Where (e) All Opposites Merge In (e&eb)The Radiance Of The Great Light Of Consciousness. The Stanzas On Vibration Teach: How Can One Who, As If Astonished, Beholds (e) His Own Nature As (=) That Which Sustains [The Existence Of Everything] Be Subject To (e&eb) This Painful Round Of Transmigration (Kusrti)? The Yogi, Recognising His True Nature To Be (=)The Supreme Subject, Is Astonished To (e) Suddenly Discover That The Individual He Thought He Was, Caught Up (e&eb) In The Trammels Of Thought And Living In A World Enmeshed In (e&eb) The Web Of Time And Space, Does Not Really Exist At All. He Experiences A 'Turning About' (Paravrtti) In (eb)The Deepest Seat Of The Doctrine Of Vibration Consciousness As He Penetrates His True Nature. The Sudden Eruption Of This Intuition (Pratibha) Arouses In Him A Cry Of Amazement As He Transcends All Thought-Constructs And, Perfectly Absorbed In His Own Nature, Is Liberated. The Path To Liberation Essentially, Spanda Doctrine Is Concerned With Two Matters. The First Is To Impart To Those Who Are Fit To Receive The Teachings A Deeper Understanding Of The Ultimate Goal Of Life (Upeya). When We Have Understood What Truly Benefits Us And Is Worth Attaining And What, On The Contrary, Is Of No Real Value But Stands In The Way Of This Attainment, We Can Begin To Make Progress Towards Our Goal. This Is Spanda Doctrine's Second Concern, Namely, To Show The Way In Which We Can Develop Spiritually Through Siva's Grace And The Right Application Of The Means To Realisation That It Teaches. When Both These Aspects Of The Teaching Have Been Correctly Understood And Applied, The Spanda Yogi Achieves A Clear And Permanent

Realisation Of His Goal And Is Liberated, Thus Fulfilling The Ultimate Aim Of The Teaching. The Doctrine Of Vibration Is Not Meant For The Spiritually Dull. It Is Not For The Worldly Whose Consciousness, Clouded By Ignorance, Is As If Dreaming, Even During The Waking State Of Daily Life, The Dream Of Its Own Thought-Constructs. The Teachings Are Meant For Those Who Are Awake (Prabuddha), Those Who, Full Of Faith And Reverence, Are Always Alert And Intent On Discerning The True Nature Of Ultimate Reality. This Reality Is Understood In Three Basic Ways. The First Is Purely Transcendental. The Stanzas Choose This Aspect As The One Which Formally Defines It Most Specifically. Ultimate Reality Transcends All The Opposites, Including Subject And Object. This Does Not Mean, However, That It Is An Unconscious Void, A Mere Absence Of All Existence. In Fact, This Negative Characterization Of Reality (Which Includes Also A Denial Of All That Is Unconscious) Implies A Positive Immanence In Which The Opposites Are United In The Oneness Of Pure Consciousness That Is Equally Siva And Spanda, His Universal Activity. These Two Seemingly Contrasting Aspects Are Reconciled In The Third, Namely, Reality Understood As The Essential Nature Of All Things. Although Universal And Everywhere The Same, It Is Understood To Be The Essential And Specific Nature Of Each Existent As Its 'Own Nature' (Svabhava).

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such

a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

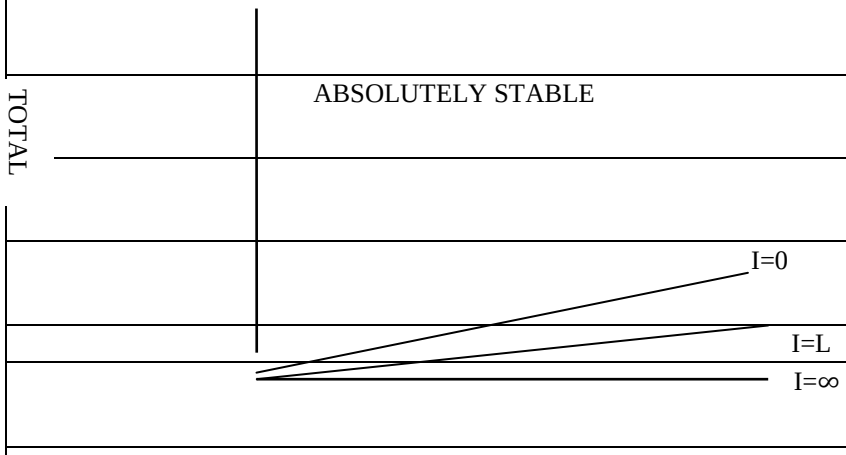
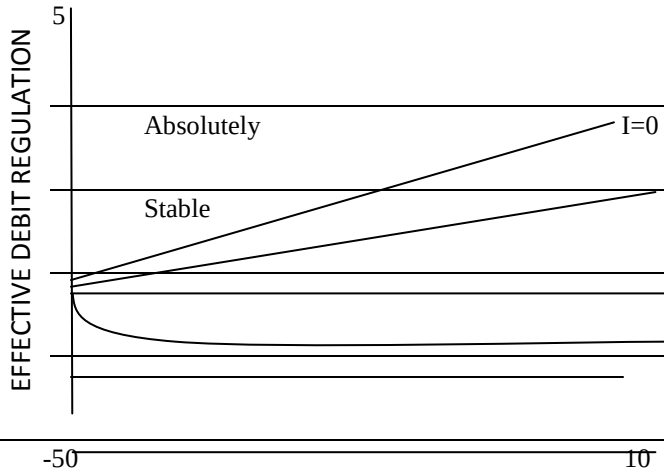



FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

This equation is of paramount importance in the following cases: According to the principles of geometric quantum mechanics, the physical characteristics of a given quantum system can be represented by geometrical features that are preferentially identified in this complex manifold. Here we construct a number of examples of such features as they arise in the state spaces for spin, spin 1, spin and spin 2 systems, and for pairs of spin systems. A study is then undertaken on the geometry of entangled states. A locally invariant measure is assigned to the degree of entanglement of a given state for a general multi-particle system, and the properties of this measure are analysed for the entangled states of a pair of spin particles. With the specification of a quantum Hamiltonian, the resulting Schrödinger trajectories induce an isometry of the Fubini–Study manifold, and hence also an isometry of each of the energy surfaces generated by level values of the expectation of the Hamiltonian. For a generic quantum evolution,

the corresponding Killing trajectory is quasiergodic on a toroidal subspace of the energy surface through the initial state. When a dynamical trajectory is lifted orthogonally to Hilbert space, it induces a geometric phase shift on the wave function. The uncertainty of an observable in a given state is the length of the gradient vector of the level surface of the expectation of the observable in that state; a fact that allows us to calculate higher orders corrections to the Heisenberg relations. A general mixed state is determined by a probability density function on the state space, for which the associated first moment is the density matrix. The advantage of a general state is in its applicability in various attempts to go beyond the standard quantum theory, some of which admit a natural phase-space characterization. **Journal of Geometry and Physics Volume 38, Issue 1, April 2001, Pages 19–53 Geometric quantum mechanics Dorje C. Brody Lane P. Hughston DOI: 10.1016/S0393-0440(00)00052-8** Expansion (unmesa) of [consciousness] or the creative intuition {pratibha) is (=) [experienced] in the interval which divides two [moments] of differentiated perception (vikalpa) It is here that they arise and disappear. The Sastras and Agamas proclaim with reasoned argument that it is (=) free of thought-constructs {Nirvikalpa) and precedes (e) Path to Liberation all mental representations of any object. None can deny that a gap exists between perceptions insofar as two moments of thought are (=) invariably divided. This [gap] is (=) the undifferentiated unity of all the countless manifestations. Similarly, in the outer more objective sphere, where change consists (e) of the alterations in the configurations of manifest appearances (Abhasa\ the transition from one to (e&eb) another corresponds to (e&eb) a phase of pure luminosity that marks (eb) the beginning of one form and the end of another

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^{\infty} \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7 \left( U_7 - V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7 \right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$then given the constraints of $f(T_7)$, $(A1.5) can\ never\ be\ satisfied\ Co$nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7 \lambda_7 e^{-T_7 \lambda_7} + [(U_7 - V_7) M_7 + Y_7] - [((U_7 -$$

V 7M 7–]**e**–7 7λ 7=0

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$\boldsymbol{Z}_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7) \cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7) \sin(Z_7 T_7) = 0$$

The warrant of Truth is ever elusive when we deconstruct the foundational justifications of our convenient notions about the way the world works. It is easy to confuse what is actually the creation of beliefs with the "discovery of Truth," a common goal of science and theology. Science offers no ontological Scientific Picture of the World but substitutes a number of applied theories as meta-theory. It's close, powerful and useful: it 'works' and allows use to do work. But the still mysterious "true" metatheory may reveal the basis of consciousness. But who knows when we will lift that veil? Permeating the living reality of our culture are certain contagious notions, fads and trends that have the ability to influence the way we think about the nature of Reality and ourselves. Some of them are toxic and can consume you; others are just whims of pop culture infectious and benign as the common cold. **Jung** described the concept of psychic contagion by certain archetypal forces inherent in the human psyche, which manifest in our spiritual lives and belief systems. He spoke of both conscious and unconscious contamination. Notions like this range from simple superstitions to scientific concepts, to urban myths. Notions sweep through our culture and insinuate themselves within its fabric, as fads, whether they are "real" or not, they can be influential. They mood alter us, make us feel we belong. An analysis of these notions is useful in distinguishing a common human phenomenon from any potential "alien influence" which may or may not be exerted on us from an unknown source. **Strange Attractors: Transference, Holography, and an Archetype Burke, J. (2003).**

**Strange Attractors: Transference, Holography, and an Archetype (Doctoral dissertation, Pacifica Graduate Institute, 2003).** In consideration to the Nobel Proze for Chemistry this year(2014), a stability nalysis of the proble in to single molexues would be a contemporaneous question that has to be taken cognisace of and given importance to . Determination of molecular structure by geometry optimization became routine only after efficient methods for calculating the first derivatives of the energy with respect to all atomic coordinates became available. Evaluation of the related second derivatives allows the prediction of vibrational frequencies if harmonic motion is estimated. More importantly, it allows for the characterization of stationary points. The frequencies are related to the eigenvalues of the Hessian matrix, which contains second derivatives. If the eigenvalues are all positive, then the frequencies are all real and the stationary point is a local minimum. If one eigenvalue is negative (i.e., an imaginary frequency), then the stationary point is a transition structure. If more than one eigenvalue is negative, then the stationary point is a more complex one, and is usually of little interest. When one of these is found, it is necessary to move the search away from it if the experimenter is looking solely for local minima and transition structures. The total energy is determined by approximate solutions of the time-dependent Schrödinger equation, usually with no relativistic terms included, and by making use of the Born–Oppenheimer approximation, which allows for the separation of electronic and nuclear motions, thereby simplifying the Schrödinger equation. This leads to the evaluation of the total energy as a sum of the electronic energy at fixed nuclei positions and the repulsion energy of the nuclei. Notable exceptions are certain approaches called direct quantum chemistry, which treat electrons and nuclei on a common footing. Density functional methods and semi-empirical methods are variants on the major theme. **Orlando Alvarez** discusses Polyakov's quantization of the string in the presence of a boundary allowing for an arbitrary topology for the world sheet. In addition to the dynamical conformal factor discovered by Polyakov, there are a finite

number of new degrees of freedom if the surface is more complicated than a sphere or a disc. The quantization of the Liouville theory in an arbitrary topology is also discussed. A one-loop calculation shows that the model is renormalizable if one performs a mass renormalization and an additive field renormalization. The renormalization group equations have a perturbative infrared unstable fixed point in all topologies. Copyright © 1983 Published by Elsevier B.V. **Nuclear Physics B Volume 216, Issue 1, 25 April 1983, Pages 125–184 Theory of strings with boundaries: Fluctuations, topology and quantum geometry Orlando Alvarez DOI: 10.1016/0550-3213(83)90490-X** Quantum geometry (the modern loop quantum gravity involving graphs and spin-networks instead of the loops) provides microscopic degrees of freedom that account for black-hole entropy. However, the procedure for state counting used in the literature contains an error and the number of the relevant horizon states is underestimated. In our paper a correct method of counting is presented by **Marcin Domagala and Jerzy Lewandowski.** Results lead to a revision of the literature of the subject. It turns out that the contribution of spins greater than 1/2 to the entropy is not negligible. Hence, the value of the Barbero–Immirzi parameter involved in the spectra of all the geometric and physical operators in this theory is different than previously derived. Also, the conjectured relation between quantum geometry and the black-hole quasi-normal modes should be understood again. **Marcin Domagala and Jerzy Lewandowski 2004 Class. Quantum Grav 21 5233 doi:10.1088/0264-9381/21/22/014 Black-hole entropy from quantum geometry** Quantum geometry predicts that a universe evolves through an inflationary phase at small volume before exiting gracefully into a standard Friedmann phase. This does not require the introduction of additional matter fields with ad hoc potentials; rather, it occurs because of a quantum gravity modification of the kinetic part of ordinary matter Hamiltonians. Authors draw the cognizance of application of the same mechanism that can explain why the present day cosmological acceleration is so tiny. DOI: http://dx.doi.org/10.1103/PhysRevLett.89.261301

**Inflation from Quantum Geometry Phys. Rev. Lett 89, 261301 – Published 12 December 2002 Martin Bojowald** The loop quantum cosmology of the closed isotropic model is studied with special emphasis on a comparison with traditional results obtained in the Wheeler-DeWitt approach by **Martin Bojowald** .Investigation includes the relation of the dynamical initial conditions to boundary conditions such as the no-boundary or the tunneling proposal and a discussion of inflation from quantum cosmology. **DOI: http://dx.doi.org/10.1103/PhysRevD.67.124023 Loop quantum cosmology, boundary proposals, and inflation Phys. Rev. D 67, 124023 – Published 23 June 2003 Martin Bojowald and Kevin Vandersloot** In loop quantum cosmology, the universe avoids a big bang singularity and undergoes an early and short super-inflation phase. During super-inflation, non-perturbative quantum corrections to the dynamics drive an inflaton field up its potential hill, thus setting the initial conditions for standard inflation. **Shinji Tsujikawa et al** show that this effect can raise the inflaton high enough to achieve sufficient e-foldings in the standard inflation era. They also analyse the cosmological perturbations generated when slow-roll is **violated after super-inflation** and show that loop quantum effects can in principle leave an **indirect signature** on the largest scales in the CMB, with some loss of power and running of the spectral index. **Shinji Tsujikawa et al 2004 Class. Quantum Grav 21 5767 doi:10.1088/0264-9381/21/24/006 Loop quantum gravity effects on inflation and the CMB Quantum geometry from quantum information** is the study of **Florian Girelli and Etera R Livine .**Loop quantum gravity defines the quantum states of space geometry as spin networks and describes their evolution in time. We reformulate spin networks in terms of harmonic oscillators and show how the holographic degrees of freedom of the theory are described as matrix models. This allows us to make a link with non-commutative geometry and to look at the issue of the semi-classical limit of loop quantum gravity from a new perspective. This work is thought of as part of a bigger project of describing quantum geometry in quantum information terms.

**Florian Girelli and Etera R Livine 2005 Class. Quantum Grav 22 3295 doi:10.1088/0264-9381/22/16/011 Reconstructing quantum geometry from quantum information: spin networks as harmonic oscillators.** Quantum geometric mechanics is the topic of study by **Dorje C. Brody Lane P. Hughston** establishing a clear relation between the quantum mechanics and quantum geometry and also the entangled states. The manifold of pure quantum states can be regarded as a complex projective space endowed with the unitary-invariant Fubini–Study metric. Note how differential gheometry concepts with Nunney's model could be entangle so that it is feasibly and viably applied to various type of model including consciousness (See Doctrine of Vibration)

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $\boldsymbol{Z}_7^2$ and $Z_7$in (A2.2):

$W_7^2 (X_7^2$-3 $M_7\ A_7\ X_7$- $M_7^2 (X_7$- $Y_7) + 2\ M_7^3\ A_7$+2 $A_7^2\ M_7^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_7$=1-cos ($Z_7\ T_7$).

In the absence of a delay the system is stable for $A_7 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7 > (X_7 - Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$A_7 > [\ X_7 - Y_7/2\ M_7$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_7, A_7) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_7 \le 2$ *is* of interest First consider the extreme values

$$Q_7(0, A_7) = A_7^2\ (\mathrm{X}_7 - Y_7)$$

$$Q_7(2, A_7) = \left(A_7 + 2\,\mathrm{M}_7\right)^2 (2A_7\,\mathrm{M}_7 - (\mathrm{X}_7 - Y_7))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7 M_7^2 - 3M_7 X_7\right) + W_7\left(3M_7 A_7^2 - 2A_7 X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7 + Y_7\right)$$

$$=2\,M_7^2 W_7^2 (M_7 + \mathrm{A}_7) + M_7\,A_7^2 W_7 + \left(2A_7 M_7 - X_7 + Y_7\right)(A_7 W_7 + M_7 W_7^2$$

$$+ (2- W_7)((2\,M_7\,X_7 + M_7\,Y_7)\,W_7 + A_7\,(X_7 + Y_7)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $\mathrm{W}_7$touches the $W$ –axis only at its minimum value in which case (A2.9) shows that

further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$ has no minimum in the range $0 \le W_7 \le 2$. Defining

$Q_7(W_7, A_7) = a_7 W_7^2 + c_7$ *where* $a_7\ b_7$ *and* $c_7$ Are functions of $A_7$ we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7 - b_7 \ge 0$$

And

$$b_7^2 - 4a_7 c_7$$

$$Case(i): Y_7 = 0$$

*Let* $Y_7 = 0$ And try the substitution A= $X_7 = /2M_7$ *in* (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$ The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}(W_7 X_7 / 2M_7)$ is positive for *W<2. Since* $Q_7$ increases with $A_7$ (see A2.9), increasing $A_7$ will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

*Case* $(ii)$: $Y_7 = y_7 X_7$ ,

$Let\ Y_7 = y_7\ X_7$And try the substitution $A_7 = X_7/(2 + A_7\ y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + n_7 y_7)$ we have

$$W_7^2\ y_7\left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2\left(2 - n_7 + n_7 y_7\right)\right]$$

$$-W_7\ X_7\left[\frac{X_7(1 + n_7 y_7)}{M_7^2(2 + n_7 y_7)^2} + 2\ y_7\left(n_7 - 1\right) + \frac{X_7^2(7 + y_7)}{M_7^2(2 + n_7 y_7)} = 0.\ (A2.12)\right.$$

Since $n_7$ and $A_7$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7), the\ condition (A2.77?) can\ be\ written\ as$

$[\frac{X_7(E_7-7)}{E_7^2}]^2[\frac{X_7^2}{M_7^2} - 4\ y_7\ E_7^2\frac{X_7}{M_7^2} + 4y_7(n_7 - 2)E_7^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$2\ E_7[y_7 E_7 + (y_7\ E_7(7 + y_7)(E_7 - n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - y_7 E_7(7 + y_7)(E_7 - n_7))^{7/2}]$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0 $W_7$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7-7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2-2y_7\frac{(E_7+n_7^2y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right)-4y_7(E_7-n_7)$$

And has the solution

$$\frac{X_7}{M_7^2}>\frac{E_7}{(E_7-7)}[y_7(E_7+n_7^2y_7)+[y_7^2(E_7+n_7^2y_7)^2+4y_7(E_7-7)(E_7-n_7)]^{7/2}]$$

Given $E_7>n_7$this lower limit to $X_7/M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7/M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$in the appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7<n_7$ $(ie.,2A_7M_7^2<X_7-Y_7$)then (A2.7)and (A2.8)show that $Q_7$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2}>2E_7[y_7E_7-(y_7E_7(7+y_7)(E_7-n_7)]^{7/2}]$$

And

$$E_7>n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2 > 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.

If $A_7^2 < 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.+ $(Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$ and $X_7$ and $Y_7$ by $W_7^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7 \lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]\, e^{2T_7\lambda_7} + M_7 V_7\, e^{2T_7\lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i} Z_7$ *and* separate

(A3.1)into real and imaginary parts

$$Z_7^2 + (M_7 + V_7) Z_7 \sin(Z_7 T_7) - (U_7 M_7 + Y_7) + (M_7(U_7 + V_7) - X_7)$$

$$\mathrm{X} \cos(Z_7 T_7) - 2M_7 V_7 \cos^2(Z_7 T_7) + M_7 V_7 = 0$$

$$(M_7 + U_7) Z_7 - (M_7 + V_7) Z_7 \cos(Z_7 T_7) + (M_7(U_7 + V_7) - X)$$

$$\mathrm{X} \sin((Z_7 T_7) - M_7 V_7 \sin(Z_7 T_7) = 0$$

Using (A.3.3) to eliminate $Z_7$ and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - cos\, Z_7\ T_7$ and $A_7 = U_7 - V_7$ we obtain

$$4\,M_7^2 V_7^2 W_7^4 + 2\,M_7 V_7\ W_7^3[(M_7 - V_7)^2 + A_7(M_7 + V_7) + 2(A_7 M_7 - X_7)]$$

$$+ W_7^2[M_7 - V_7)^2(2\,A_7\ M_7 + Y_7 - X_7) + (2A_7\ M_7 - X_7)^2 - A_7(M_7 - V_7)(2A_7\ M_7 - X_7)$$

$$+ M_7 V_7\,(A_7^2 + 4(X_7 + Y_7))]$$

$$+ W_7[A_7^2 + 2X_7)(2\,A_7\ M_7 - X_7) + 2A_7 Y_7(M_7 + V_7) - M_7 A_7^3] + A_7^2(X_7\ M_7 - Y_7) = 0$$

Being a quartic in $W_7$ (A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7 > 0$$

$$and$$

$$C_7 > (X_7 - Y_7)/2\,M_7$$

*recalling that* $A_7 = U_7$ and defining $C_7 = U_7 + V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0, U_7 V_7 = A_7^2(X_7 + Y_7)$$

$$Q_7\left(2, U_7 V_7\right) = \left(C_7 + 2M_7^2\right)^2 (2C_7 M_7 + Y_7 - X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0$W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

## Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

**Module numbered five**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This is the condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference

between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \quad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t - T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\left(R(t)\right)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

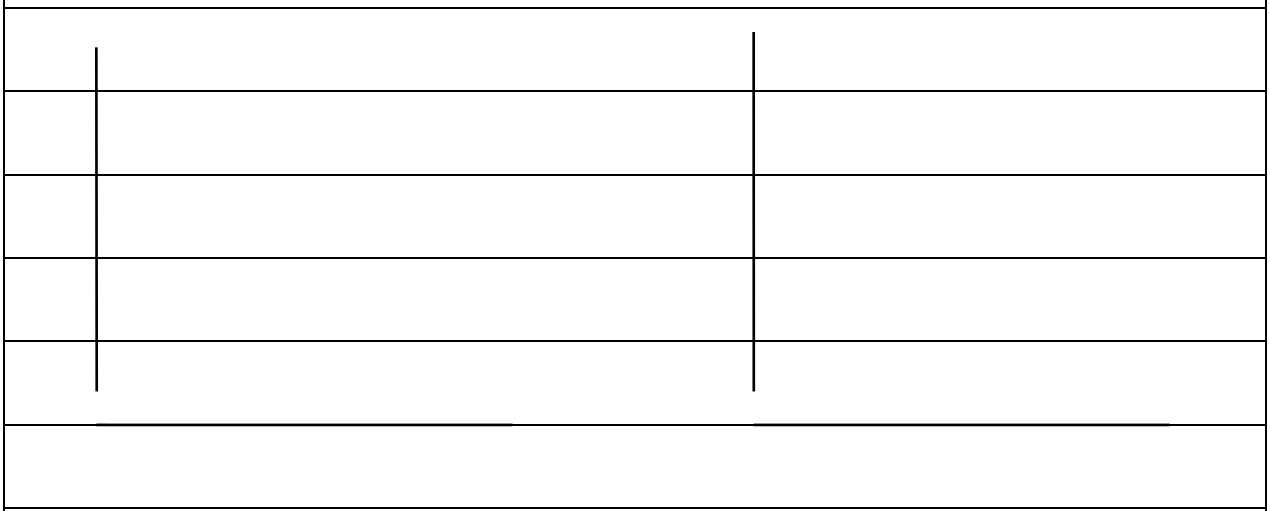
Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time

delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the



TIME DELAY

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period,

which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

As it is usually expressed outside of phenomenological texts, to undergo a conscious experience necessarily means that there is something it is like for the subject to have that experience (Nagel 1974; Searle 1992). This is obviously true of bodily sensations like pain. But it is also the case for perceptual experiences, experiences of desiring, feeling, and thinking. There is something it is like to taste chocolate, and this is different from what it is like to remember what it is like to taste chocolate, or to smell vanilla, to run, to stand still, to feel envious, nervous, depressed or happy, or to entertain an abstract belief. Yet, at the same time, as I live through these differences, there is something experiential that is, in some sense, the same, namely, their distinct first-personal character. All the experiences are characterized by a quality of mineness or for-me-ness, the fact that it is I who am having these experiences. All the experiences are given (at least tacitly) as my experiences, as experiences I am undergoing or living through. All of

this suggests that first-person experience presents me with an immediate and non-observational access to myself, and that consequently (phenomenal) consciousness consequently entails a (minimal) form of self-consciousness. To put it differently, unless a mental process is pre-reflectively self-conscious there will be nothing it is like to undergo the process, and it therefore cannot be a phenomenally conscious process. The mineness in question is not a quality like being scarlet, sour or soft. It doesn't refer to a specific experiential content, to a specific what; nor does it refer to the diachronic or synchronic sum of such content, or to some other relation that might obtain between the contents in question. Rather, it refers to the distinct givenness or the how it feels of experience. It refers to the first-personal presence or character of experience. It refers to the fact that the experiences I am living through are given differently (but not necessarily better) to me than to anybody else. It could consequently be claimed that anybody who denies the for-me-ness of experience simply fails to recognize an essential constitutive aspect of experience. Such a denial would be tantamount to a denial of the first-person perspective. It would entail the view that my own mind is either not given to me at all — I would be mind- or self-blind — or is presented to me in exactly the same way as the minds of others. There are also lines of argumentation in contemporary analytical philosophy of mind that is close to and consistent with the phenomenological conception of pre-reflective self-awareness. Alvin Goldman provides an example: [Consider] the case of thinking about x or attending to x. In the process of thinking about x there is already an implicit awareness that one is thinking about x. There is no need for reflection here, for taking a step back from thinking about x in order to examine it…When we are thinking about x, the mind is focused on x, not on our thinking of x. Nevertheless, the process of thinking about x carries with it a non-reflective self-awareness (Goldman 1970, 96). A similar view has been defended by Owen Flanagan, who not only argues that consciousness involves self-consciousness in the weak sense that there is something it is like for the subject to have the experience, but also speaks of the low-level self-consciousness involved in experiencing my experiences as mine (Flanagan 1992, 194). As Flanagan quite correctly points out, this primary type of self-consciousness should not be confused with the

much stronger notion of self-consciousness that is in play when we are thinking about our own narrative self. The latter form of reflective self-consciousness presupposes both conceptual knowledge and narrative competence. It requires maturation and socialization, and the ability to access and issue reports about the states, traits, dispositions that make one the person one is. To claim that every kind of self-consciousness is conceptual is overly cognitive. Bermúdez (1998), to mention one further philosopher in the analytic tradition, argues that there are a variety of **nonconceptual forms of self-consciousness** that are "logically and ontogenetically more primitive than the higher forms of self-consciousness that are usually the focus of philosophical debate" (1998, 274; also see Poellner 2003). This growing consensus across philosophical studies supports the phenomenological view of pre-reflective self-consciousness. (References: Stanford Encyclopedia, Post Modersnist Theories, Internet Encyclopedia of Philosophy) Korzybski pointed out that there are two ways to slice easily through life; to believe everything or to doubt everything. Both ways save us from thinking. So, from time to time you have to give yourself a Virtual Reality Check. The hallmark of postmodern philosophy has been disbelief or skepticism of all "metanarratives," or translations of reality. Postmodernism has even turned its profound skepticism on such important humanist concepts as "objective truth" and reason.Yet, for a deconstructionist postmodern society, individually we are still riddled with superstition and gullibility, and open to manipulation through our belief systems as any politician, philosopher, clergy, or salesperson will attest. Further, most people are painfully naive when it comes to even the simplest scientific understanding. Most of us don't have a clue about the fundamental nature of physical reality or our own psychological nature, and our ability to be fooled by our senses and mind, or methods of social and technological persuasion. We may think we know, but the penetration is deeper than we dare imagine. Nature rejects the naiveté that seeks absolute truth. We are beginning to realize, individually and culturally, that "realities" are all human constructions. The task becomes one of "catching ourselves in the act" of creating our own "reality" from the flow of events. Human truth is always an engagement of mind with experience.. We don't need to fear the collapse of our

personalistic belief system (the "box" we live in), nor our belief in absolute truth. **A strong desire to engage in the "quest for uncertainty" complements our anxiety that perhaps there is no absolute, objective ground to reality**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t-T_8)F_8(R_8(t-T_8) - N_8(t)M_8((R_8(t) \qquad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8} + \frac{N_8 dG_8}{dR_8} - \frac{dB_8}{dR_8} > \frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8} + \frac{dM_8}{dR_8}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_8 = \frac{N_8 G_8}{2M_8}\frac{d(8nF_8)}{dR_8} (\geq 0) \quad (8a)$$

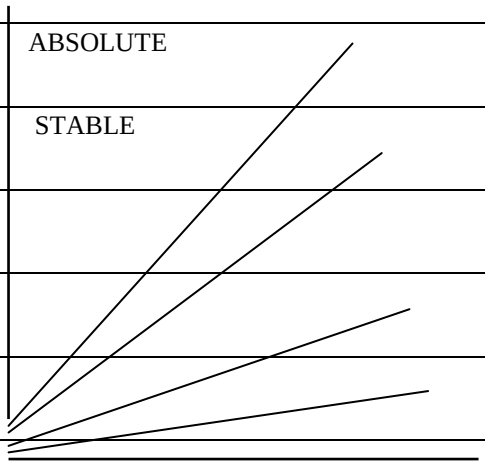
$$I_8 = -\frac{N_8 G_8}{2M_8}\frac{d(8nM_8}{dR_8} (\geq 0) \quad (8b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8 > L_8 - I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_8 > L_8 - I_8$, when $P_8^2/4 > I$ (8a)

$P_8 > L_8 - I_8 + \frac{(I_8 - P_8^2}{(L_8 - I)}$ when $P_8^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

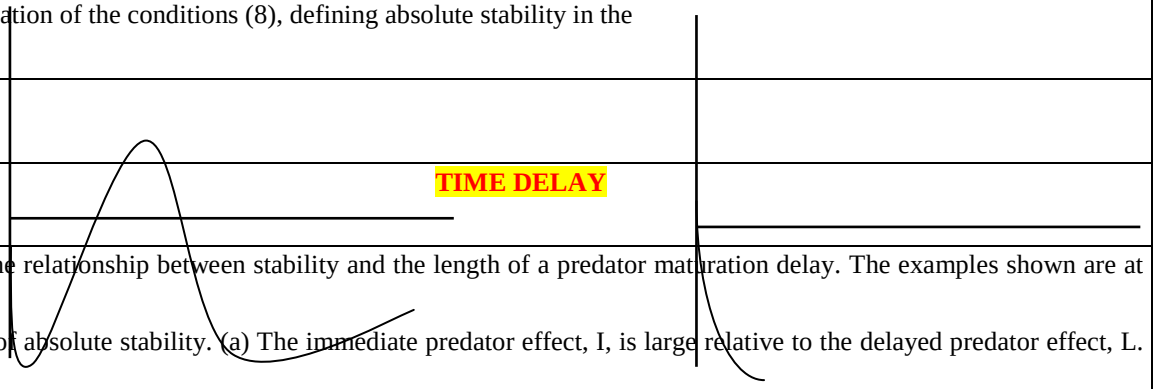



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this

rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior , we must note that the model could be applied to psychological systems too, albeit we are still not sure of the fact how the wave function collapses thereby due to perception.

Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility. He privatizes the sublime, having learned from the fate of his predecessors that the public can never be more than beautiful."(11) Notice is made here of what is indeed the telling feature. Levinas is praised, for example, for his alarm when "the social will [is] sought in an ideal of fusion . . . the subject losing himself in a collective representation, in a common ideal. . . . It is the collectivity which says 'us,' and which, turned toward the intelligible sun, toward the truth, experience, the other at his side and not face to face with him. The tellingly titled D'un ton apocalyptique contains passages that are particularly excited: These people situate themselves outside the ordinary, but they have in common this: they describe themselves as having an immediate and intuitive relation with mystery. And they want to attract, to seduce, and lead others to the mystery, through mystery. This agogic function of the leader of men, of the duce, of the Führer, of the leader, places him above the crowd that he manipulates with the aid of a small numbers of adepts joined together in a sect with a secret language, a clique or a small party with its

ritualized practices. The mystifiers pretend to have exclusive access to the privilege of a secret mystery. . . . The revelation or the unveiling of the secret is something that they jealously reserve for themselves. Jealousy is here a major characteristic. There is not narcissism and non-narcissism; there are narcissisms that are more or less comprehensive, generous, open, extended. What is called non-narcissism is in general but the economy of a much more welcoming, hospitable narcissism, one that is much more open to the experience of the other as other. I believe that without a movement of narcissistic reappropriation, the relation to the other would be absolutely destroyed; it would be destroyed in advance. The relation to the other--even if it remains asymmetrical, open, without reapporpriation--must trace a movement of reappropriation in the image of oneself for love to be possible, for example. Love is narcissistic. Beyond that, there are little narcissisms, there are big narcissisms. . . .(23) **Justice of the Pieces Deconstruction as a Social Psychology Douglas Collins** Sabine Wilke cogently argues for the rise of the aesthetic dimension in philosophical writing by showing how the metaphorical and poetic style of Heidegger, the principles of phrase configurations with Adorno, or the topographic writing of derrida replace through their very form something what traditionally would have been subject to discursive exposure The concern shifts from telling to showing, showing first and foremost that something remains untold. An exemplary book in this aesthetic drive is Avital Ronell's Telephone Book where the playful layout and explicit break with traditional design interrupt the sequential logic of the paragraphs thereby performing the very argument at stake: open up new circuits of meaning and signification. The engagement with form and publishing formats is only consequential since media theory as exposed by Walter J. Ong, Eric A. Havelock, Parry Miman or Marshall McLuhan besides many others made it clear that different communication technologies condition the mode of thought of a given culture The fractals of Benoit Mandelbrot or the butterfly shaped Lorenz attractor are the popular icons of this turn in the sciences. For communication studies and philosophy, Vilém Flusser was one of the pioneer thinkers reflecting the rise of what he called “techno-images.” The photograph is the paradigmatic case which condenses the whole crises of Western

thought and culture on its grainy surface. It subverts the linear logic of writing.6 Flusser's central diagnosis detects a shift from the grand narratives associated with the literate mind and Modernity towards creating and envisioning the world through image based practices. Writing himself in a succinct, visually inspired style, Flusser contents that the web of meaning, spun throughout the centuries by our textual practices, has become transparent. We discover the void between the letters through which the world escapes: "We are alienated from the world of text, precisely because we see through them as our own product."7 As a consequence we are confronted with a world of dissociated facts and information bits. Flusser hails this fragmented world as the advent of the zero dimensional point-universe that will provide the new raw material for our image based practices. Techno-images are nothing else than a conscious, inherently artistic drawing-together of point-elements as exemplified first and foremost by the silver-grains of the photograph but also by more recent pixel based techno-images. Similar to blueprints, designs or statistical curves, the weather map for example visualizes and contracts a tremendous amount of meteorological data from widely distributed stations into a single image of high- and low pressure

areas, temperatures, wind speeds, etc. Techno-images "compute" in this sense the universe of unimaginable, abstract information-bits into a concrete surface – the image. Echoing Paul Feyerabend, these image practices have to be understood as creative practices where the truth of scientific models has been replaced by its aesthetic quality **Aesthetics, Writing, Networked Computers Jörg Muller Dissertation Submitted to the Division of Media and Communications of The European Graduate School in Candidacy for the Degree of Doctor of Philosophy**

**Lothar Schäfer (1) and Sisir Roy (2) in** Quantum Reality, the Importance of Consciousness in the Universe, the Discovery of a Non-Empirical Realm of Physical Reality, and the Convergence with Ancient Traditions of Indian and Western Philosophy investigate metaphysical aspects that include 1) the discovery of a non-empirical part of physical reality in a realm of potentiality; 2) the emanation of the empirical world out of a realm of non-material forms; 3) the discovery that the nature of physical reality is that of an indivisible Wholeness – the One; and 4)

panpsychism: the possibility that the One is aware of its processes like a Cosmic Consciousness. The convergence of powerful traditions of seemingly disparate cultures is particularly important to point out in the present process of globalization, when a unifying view is needed to avoid controversy and conflict. To quote them in extenso : This is the nonlocality of the quantum world. Menas Kafatos and Robert Nadeau (1990) have drawn a remarkable conclusion from this phenomenon: If reality is nonlocal, the nature of the universe is that of an undivided wholeness. Because our consciousness has emerged from this wholeness and is part of it, it is possible to conclude that an element of consciousness is active in the universe: a Cosmic Consciousness. Pre-reflective self-consciousness is pre-reflective in the sense that (1) it is an awareness we have before we do any reflecting on our experience; (2) it is an implicit and first-order awareness rather than an explicit or higher-order form of self-consciousness. Indeed, an explicit reflective self-consciousness is possible only because there is a pre-reflective self-awareness that is an on-going and more primary self-consciousness. Although phenomenologists do not always agree on important questions about method, focus, or even whether there is an ego or self, they are in close to unanimous agreement about the idea that the experiential dimension always involves such an implicit pre-reflective self-awareness. In line with Edmund Husserl (1959, 189, 412), who maintains that consciousness always involves a self-appearance (Für-sich-selbst-erscheinens), and in agreement with Michel Henry (1963, 1965), who notes that experience is always self-manifesting, and with Maurice Merleau-Ponty who states that consciousness is always given to itself and that the word ‘consciousness’ has no meaning independently of this self-givenness (Merleau-Ponty 1945, 488), Jean-Paul Sartre writes that pre-reflective self-consciousness is not simply a quality added to the experience, an accessory; rather, it constitutes the very mode of being of the experience: This self-consciousness we ought to consider not as a new consciousness, but as the only mode of existence which is possible for a consciousness of something (Sartre 1943, 20 [1956, liv]). The notion of pre-reflective self-awareness is related to the idea that

experiences have a subjective 'feel' to them, a certain (phenomenal) quality of 'what it is like' or what it 'feels' like to have them.

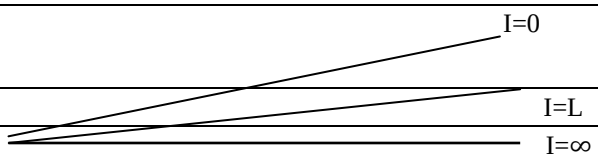
Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

TOTAL

ABSOLUTELY STABLE

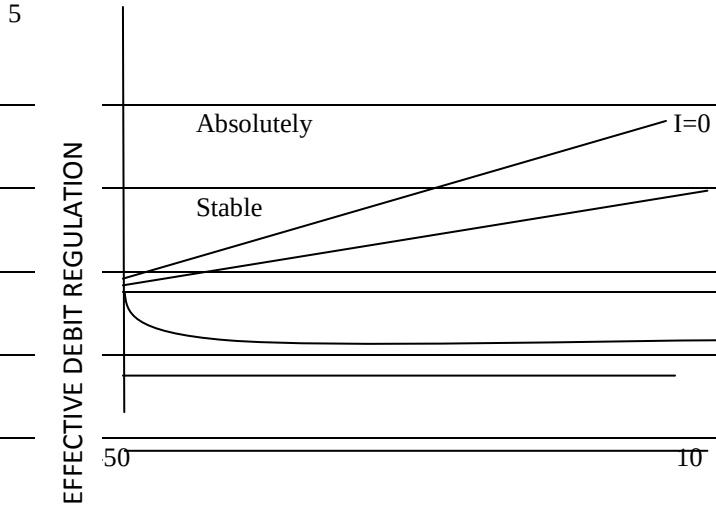



**DELAYED DEBIT EFFECT**

FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

## Utility Analysis And neural Network

Caroline Williams investigated affective density of the political and its effect on our understanding of political subjectivity. Taking up Spinoza's challenge to think about affect beyond corporeal embodiment, he argues that there is a modality of affectivity that cannot simply be inscribed within the borders of subjectivity. theorising affect as an impersonal force anchored in a relational ontology that gives due recognition to the circulation of affects, as well as to their ambivalent structure in creating sites of identification, and he utilises this ontology to reflect on the dynamic of the political and the shape of political subjectivity. Williams argues that Spinoza's philosophy (through ideas of conatus and imagination) offers the conceptual resources to reconfigure the composition of affective subjectivity as a transindividual social bond and as an unconscious dynamic of ethico-political existence. **Subjectivity (2010) 3, 245–262. doi:10.1057/sub.2010.15 Affective processes without a subject: Rethinking the relation between subjectivity and affect with Spinoza Caroline Williams** Phenomenological Time bears ample evidence and infallible observatory and apodictic evidence albeit with a sense of chagrined circumvention and defeated disgruntlment circularity of both psychological and phenomenological interpretations of art which are sometimes insightful, but most often work on the logic of cyclical time and unavoidable drives and representations played out as a ‘natural’ order: where ‘men can’t avoid being boys’ a circularity which valorizes a helplessness/truthfulness conundrum. One begins with a classical phenomenological fallacy with Duchamp’s Fountain, does the urinal keep returning or do we keep returning to the urinal? Put in this way, it is a male question, something that phenomenology conveniently conceals when positing the “I”, or indeed, the “we”. But we must be careful not to simply assume that the phenomenological first person is white, heterosexual, able-bodied and male, providing a foundation upon which knowledge should be based as a universal condition. Deleuze’s second critique of time is linked to Kant. This model of time extends not only to phenomenology but to the phenomenological interpretation

of art and is seen as a straight line. In the Critique of Pure Reason, Kant cuts up circular time and reassembles it as a series of sensory experiences which memory processes and which constitute subjectivity. Whereas the circular model of time (and interpretation) inscribes fatalism into the subject and the world, the phenomenological method is in danger of solipsism, where the subject constitutes the world. Much of embodied philosophy and enactivist theories of cognition go down this line with a plethora of art historical analyses by the likes of Rosalind Krauss, Michael Fried and others taking the lead from Merleau-Ponty. These approaches successfully challenge the **homunculus model of a Cartesian mind-body** dualism where there is a supervisor in the mind who processes experience and even consciousness and to whom all representations are directed. But phenomenology is mainly concerned with perceptual consciousness, the processing of colors, shapes, forms, motion in art and the world. The senses and sensorimotor processes are the fundamental pathways by which the interior is connected to the exterior in a chiasm of self-constitution in the world. The problem with this is that we explore the art through our sense perceptions, and what does art reveal to us?—our sense perceptions! Such solipsistic approaches have real problems relating complex and abstract conceptual production to bodily and sensorimotor contingencies, and they nearly always posit the subject as the site and object of discovery in the world. To be fair, phenomenology has moved on to rather more sophisticated models of intersubjectivity, taking into account the specificities of different kinds of bodies and ways of being beyond the white, male heterosexual as the common denominator, yet many naïve approaches to art perpetuate the notion that perceptual, rather than conceptual processes and evolutionary neural patterns are primary, fundamental ways to understand the art experience In fact, these aspects are only part of art production and reception, aspects which are often filtered out or made known to us in order for us to parody such trigger responses, as in much of Duchamp's work, which holds in contempt such retinal approaches to art. If we take the **subject of abstract expressionism** in the 1950s or later minimalist art, art history, teaching and interpretation still insists on putting the act of perceiving as fundamental to the constitution not only of the phenomenological first

person but also of art itself. The male act of usually standing up, which consists of directing a stream with the hand(s) into a container or on a flat surface, sand or snow in order to leave patterns and traces (the usual child's prank) underlies the logic of macho 1950s Abstract Expressionist culture, especially that of Jackson Pollock, so effectively parodied by Carolee Schneeman's Vagina Scroll, 1975. Here, painting or writing is a kind of a discharge which invests the **object and space with** the phenomenological "I", disciplined by existential struggle and reflected back by the traces, signatures and gestures left behind in the work of art, rediscoverable as an element in the making of other subjectivities.It is precisely this kind of solipsism which is blind to the art event as something well in excess of re-establishing habits of subjective self-identity through the production and interpretation of art. The attempt to reinforce a phenomenological identity upon the event by personalizing sense experience is unable to see art as something which remodels the phenomenological 'I' and neural plasticity itself. We learn new technologies and redraw our neural pathways through such new and spontaneous experiences which draw us into intersubjective worlds. Duchamp's Fountain and Deleuze's Repetition and Différence Grégory Minis sale (I have made some changes) The relation of different groups to deconstruction is settled by its relation to groups. An American enthusiast--one who has had a role in the development of our affection for Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^\infty \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$then given the constraints of $f(T_8)$, $(A1.5) can\ never\ be\ satisfied\ Co$nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8 \lambda_8 e^{-T_8 \lambda_8} + [(U_8 - V_8) M_8 + Y_8] - [((U_8 - V_8) M_8 -] e^{-T_8 \lambda_8} = 0$$

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$\boldsymbol{Z}_8^2 + M_8 Z_8 \sin(Z_8 T_8) - (A_8 M_8 + Y_8) + (A_8 M_8 + X_8) \cos(Z_8 T_8) = 0$$

$$(A_8 M_8) Z_8 - M_8 Z_8 \cos(Z_8 T_8) + (A_8 M_8 - X_8) \sin(Z_8 T_8) = 0$$

Where $A_8 = U_8 - V_8$ Using (A2.3), substitute for $\boldsymbol{Z}_8^2$ and $Z_8$ in (A2.2):

$W_8^2 (X_8^2 - 3 M_8 A_8 X_8 - M_8^2 (X_8 - Y_8) + 2 M_8^3 A_8 + 2 A_8^2 M_8^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_8 = 1 - \cos(Z_8 T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8 / 2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_8 \le 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2 (X_8 - Y_8)$$

$$Q_8(2, A_8) = (A_8 + 2M_8)^2 (2A_8 M_8 - (X_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2 \left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8 \left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8 \left(2X_8 + Y_8\right)\right)$$

$$+2A_8 \left(X_8 + Y_8\right)$$

$$= 2M_8^2 W_8^2 (M_8 + A_8) + M_8 A_8^2 W_8 + (2A_8 M_8 - X_8 + Y_8)(A_8 W_8 + M_8 W_8^2$$

$$+ (2 - W_8)((2M_8 X_8 + M_8 Y_8) W_8 + A_8 (X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $W_8$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in *A* will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range $0 \le W_8 \le 2$.Defining

$Q_8(W_8, A_8) = a_8 W_8^2 + c_8$ *where* $a_8$ $b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \square - b_8 \square 0$$

And

$$b_8^2 \square 4a_8 c_8$$

$$Case(i): Y_8 = 0$$

$Let\ Y_8 = 0$ And try the substitution A= $X_8 = /2M_8\, in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8 /$ $2M_8$) is positive for *W<2. Since* $Q_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\ X_8$ ,

$Let\ Y_8 = \ y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 + \ n_8 y_8)$ we have

$$W_8^2\, y_8 \left[X_8 n_8 \frac{(1 + n_8 y_8)}{(2 + n_8 y_8)} + M_8^2 \left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8 X_8 \left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2}+2y_8\left(n_8-1\right)+\frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)}\right]=0.\ (A2.12)$$

Since $n_8$ and $A_8$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8=(2+n_8 y_8)$, *the condition*$(A2.88?)$*can be written as*

$$\left[\frac{X_8(E_8-8)}{E_8^2}\right]^2\left[\frac{X_8^2}{M_8^2}-4y_8 E_8^2\frac{X_8}{M_8^2}+4y_8(n_8-2)E_8^3\right]$$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8/W_8^2)$ can be expressed as

$$2E_8\left[y_8E_8+(y_8E_8(8+y_8)(E_8-n_8))^{8/2}\right]>\frac{X_8}{M_8^2}>2E_8\left[y_8E_8-y_8E_8(8+y_8)(E_8-n_8))^{8/2}\right]$$

Thoughts on the Social and Political Implications of Correlationism .Lion Mirror Truth be told, as my thought has evolved the issue of correlationism had fallen off the radar for me. Somehow the debate had come to seem too "philosophical" to me, too "scholastic", too remote from what interests me: understanding why social assemblages are organized as they are, how power functions in social assemblages, and what we might do to address that power and change things. Somehow the question of whether or not we can get out of the correlation between thinking and being just came to seem remote from these sorts of issues. Somehow it seemed too epistemological. ASIDE: Numerous discussions over the years have led me to believe that the debate over correlationism is poorly understood (or maybe I just don't understand it). On countless occasions I've heard people say "of course we must relate to things in order to know them." Well yeah, of course! I don't think this is what the critic of correlationism is getting at. It seems to me that correlationism is something more robust than the theses that we must relate to something to know it.

Correlationism instead seems to require the theses that thought and being are indiscernible. Put more concretely, the correlationist is someone who argues that we either a) can never tell whether being is merely a construction of our thought (weak correlationism), or b) who argues that thought actually constructs being (strong correlationism). In other words, correlationism is another name for idealism. One can hold that we must relate to something in order to know it without being a correlationist. As an aside I should also add that I am a correlationist about some things. For example, I think money is something constructed by society and is therefore a strong correlationist when it comes to money. At any rate, for a long time I'd become rather indifferent to debates about correlationism and philosophies of access. I had learned the lessons of speculative realism– which I could have also learned, I think, from Deleuze and Guattari, the new materialist feminists, actor-network theorists such as Latour –and had moved on. However, occasionally you come across a tone of phrase that pitches something in a different light. In The Cut of the Real, Katerina Kolozova writes, …the political problem of contemporary philosophy identified by the 'new realists' is, in fact, the product of a more fundamental epistemic problem. In his book After Finitude, Quentin Meillassoux calls this problem 'correlationism' and identifies it as an essentially post-Kantian legacy, which continues to dominate and limit philosophy. As a matter of fact, correlationism lies at (eb) the heart of postmodern theory and consists in the premise that thought can only 'think itself,' that the real is inaccessible to knowledge and human subjectivity, and that there is nothing but discursive constructs that fully determine thinking and that are meth0dologically accounted for all the way down. (1 – 2) Thought thinking only itself: Thought only encountering itself. In the jargon of postmodern and poststructuralist lingo, this would be the thesis of infinite semiosis, where signs ("thoughts") only ever relate to other signs. Within this framework, discursivity comes to be the hegemonic framework defining all of being. At the level of politics and social theory more generally, if the correlationist thesis is true the

consequences are clear: all social phenomena are discursive and all solutions to social and political problems will be discursive. The sole sphere of the political will be the discursive and all questions of politics will be questions of speech-acts and interpretation. The problem here is not that many theorists recognize that the discursive and semiotic plays an important role in the social and the political. It does and I've repeated this tirelessly. The problem is with what happens when thought or the semiotic becomes a hegemon, an "all", foreclosing our ability to recognize other forms of power. What I've wanted to say is that not all power functions discursively. In my last post and elsewhere I spoke of some other forms of politics: Thermopolitics: The politics surrounding **energy in the form of calories and fuels such as gasoline** and coal, and how our life and our very bodies are structured by energy dependencies and by being trapped in particular distributive networks that render these forms of energy available. I'm being quite literal when I speak of energy, talking about the effects, for example, of the absence of food in certain educational environments on cognition, for example; and am generally hostile to metaphorical extensions of the concept of energy which I see as erasing the dimension of real materiality. Geopolitics: The role that features of natural and built geography such as mountain ranges, rivers, oceans, soil conditions, roads, housing design, etc., play in the form that social relations take and how they impact individual bodies. Chronopolitics: The way in which the structurations of time organize what is possible for us. For example, the structuration of the working day, how much we can say and comprehend at any given time, the impact of things like the invention of the clock, etc. **Oikopolitics: This would be the domain of political economy described so well by Marxists**. So five different types of politics: Semiopolitics (or what currently dominates critical theory), thermopolitics, geopolitics, chronopolitics, and oikopolitics. No doubt there are other sites of the political or political struggle that we could speak of, but this is a good start. Also, it should be obvious that these aren't exclusive domains, but are entangled in all sorts of important

ways. For example, something might take place at the level of semiopolitics (speech, law, rhetoric, norms, and communication) that has all sorts of impact at the level of thermopolitics. Congress might decide to cut programs that fund school meal programs. This, in turn, will have a thermodynamic impact on those students that go without the calories they need developmentally and cognitively to function in a particular way. There is an entanglement here of semiopolitical and thermopolitical domains. The young student here has been constrained both at the level of **semiotic phenomena and thermodynamic structures**. The point is that if true, semiotic intervention (speech-acts, protests, interpretations, deconstructions, etc) will not be an appropriate response to all political problems because social formations are not entirely structured by the semiotic. The child in that school does not suffer from a lack of the right signs, but from a lack of calories needed to run the engine of his thought and body. Certainly semiotic interventions might be needed to render that energy available, but it is the energy itself that is at issue and the absence of that energy that forms the spider web entangling him in his position. A correlationist perspective tends to erase this as even being a site of the political. **Larval Subjects January 28, 2014 Ontological Anarché Posted by larval subjects under Uncategorized January 27, 2014 Roden on Pluralism Over at enemy industry Roden has an excellent post up on the pluralism discussion. Check it out here. January 25, 2014 Different Senses of Pluralism and Ontology Posted by larval subjects under Uncategorized Responding to one of my comments over at his blog, Jeremy Trombley:** I have deleted restricted, reformulated some sentences .Kindly bear with me. Intention is to make the discursive enucleation consistent with the subject in question**.** Recourse to epistemology of the quantum mechanics generic and fervor of string theory is a sine qua non for the further comprehension of the theory and ontology .How would the world appear to us if its ontology was that of classical mechanics but every agent faced a restriction on how much they could come to know about the classical state? We show that in most respects

it would appear to us as quantum. The statistical theory of classical mechanics, which specifies how probability distributions over **phase space evolve** under Hamiltonian evolution and under measurements, is typically called Liouville mechanics, so the theory we explore here is Liouville mechanics with an epistemic restriction. The particular epistemic restriction we posit as our foundational postulate specifies two constraints. The first constraint is a classical analog of Heisenberg's uncertainty principle; the second-order moments of position and momentum defined by the phase-space distribution that characterizes an agent's knowledge are required to satisfy the same constraints as are satisfied by the moments of position and momentum observables for a quantum state. The second constraint is that the distribution should have maximal entropy for the given moments. Starting from this postulate, authors derive the allowed preparations, measurements, and transformations and demonstrate that they are isomorphic to those allowed in Gaussian quantum mechanics and generate the same experimental statistics. They argue that this reconstruction of Gaussian quantum mechanics constitutes additional evidence in favor of a research program wherein quantum states are interpreted as states of incomplete knowledge and that the phenomena that do not arise in Gaussian quantum mechanics provide the best clues for how one might reconstruct the full quantum theory. I: http://dx.doi.org/10.1103/PhysRevA.86.012103 **Reconstruction of Gaussian quantum mechanics from Liouville mechanics with an epistemic restriction Phys. Rev. A 86, 012103 – Published 10 July 2012 Stephen D. Bartlett, Terry Rudolph, and Robert W. Spekkens.** Deleuze in fact characterizes modes of existence, with their powers and capacities. The answer is this: Deleuze approaches modes of existence, ethically speaking, not in terms of their will, or their conscious decision making power (as in Kant), nor in terms of their interests (as in Marx, for example), but rather in terms of their drives. For Deleuze, conscious will and preconscious interest are both subsequent to our unconscious drives, and it is at the level of the drives that we have to aim our ethical analysis. Here, I would like to

focus on two sets of texts on the drives taken, not from Nietzsche and Spinoza, but rather from Nietzsche and Leibniz (Leibniz being one of the first philosophers in the history of philosophy to have developed a theory of the unconscious). The first set of texts comes from Nietzsche's great early book entitled Daybreak, published in July 1881. Nietzsche first approaches the question of the drives by giving us an everyday scenario: "Suppose we were in the market place one day," he writes, "and we noticed someone laughing at us as we went by: this event will signify this or that to us according to whether this or that drive happens at that moment to be at its height in us—and it will be a quite different event according to the kind of person we are. One person will absorb it like a drop of rain, another will shake it from him like an insect, another will try to pick a quarrel, another will examine his clothing to see if there is anything about it that might give rise to laughter, another will be led to reflect on the nature of laughter as such, another will be glad to have involuntarily augmented the amount of cheerfulness and sunshine in the world—and in each case, a drive has gratified itself, whether it be the drive to annoyance, or to combativeness or to reflection or to benevolence. This drive seized the event as its prey. Why precisely this one? Because, thirsty and hungry, it was lying in wait" (D 119). This is the source of Nietzsche's doctrine of perspectivism ("there are no facts, only interpretations"), but what is often overlooked is that, for Nietzsche, it is our drives that interpret the world, that are perspectival—and not our egos, not our conscious opinions. It is not so much that I have a different perspective on the world than you; it is rather that each of us has multiple perspectives on the world because of the multiplicity of our drives—drives that are often contradictory among themselves. "Within ourselves," Nietzsche writes, "we can be egoistic or altruistic, hard-hearted, magnanimous, just, lenient, insincere, can cause pain or give pleasure" (Parkes, pp. 291-292). We all contain such "a vast confusion of contradictory drives" (WP 259) that we are, as Nietzsche liked to say, multiplicities, and not unities. Moreover, these drives are in a constant struggle or combat with each other:

my drive to smoke and get my nicotine rush is in combat with (but also coexistent with) my drive to quit. This is where Nietzsche first developed his concept of the will to power—at the level of the drives. "Every drive is a kind of lust to rule," he writes, "each one has its perspective that it would like to compel all the other drives to accept as a norm" (WP 481). DELEUZE AND THE QUESTION OF DESIRE: TOWARD AN IMMANENT THEORY OF ETHICS Daniel W. Smith. **In on a Possible Physical Metatheory of Consciousness Miroljub Dugic et al** show that the modern quantum mechanics, and particularly the theory of decoherence, allows formulating a sort of a physical metatheory of consciousness. Particularly, the analysis of the necessary conditions for the occurrence of decoherence, along with the hypothesis that consciousness bears (more-or-less) well definable physical origin, leads to a wider physical picture naturally involving consciousness. This can be considered as a sort of a psycho-physical parallelism, but on very wide scales bearing some cosmological relevance. Open Systems & Information Dynamics Volume 9, Number 2, 153 (2002) arXiv: quant-ph/0212128. On a broader perspective and background, the theoretical computation of the universe is justified.IN physics and cosmology, digital physics is a collection of theoretical perspectives based on the premise that the universe is, at heart, describable by information, and is therefore computable. Therefore, according to this theory, the universe can be conceived of as either the output of a deterministic or probabilistic computer program, a vast, digital computation device, or mathematically isomorphic to such a device. Digital physics is grounded in one or more of the following hypotheses; listed in order of decreasing strength. The universe or reality: is essentially informational (although not every informational ontology needs to be digital); is essentially computable (the pancomputationalist position); can be described digitally; is in essence digital; is itself a computer (pancomputationalism); is the output of a simulated reality exercise. Following Jaynes and Weizsäcker, the physicist John Archibald Wheeler wrote the following: [...] it is not unreasonable to imagine that

information sits at the core of physics, just as it sits at the core of a computer. (John Archibald Wheeler 1998: 340) It from bit. Otherwise put, every 'it'—every particle, every field of force, even the space-time continuum itself—derives its function, its meaning, its very existence entirely—even if in some contexts indirectly—from the apparatus-elicited answers to yes-or-no questions, binary choices, bits. 'It from bit' symbolizes the idea that every item of the physical world has at bottom—a very deep bottom, in most instances—an immaterial source and explanation; that which we call reality arises in the last analysis from the posing of yes–no questions and the registering of equipment-evoked responses; in short, that all things physical are information-theoretic in origin and that this is a participatory universe. (John Archibald Wheeler 1990: 5) David Chalmers of the Australian National University summarised Wheeler's views as follows: Wheeler (1990) has suggested that information is fundamental to the physics of the universe. According to this 'it from bit' doctrine, the laws of physics can be cast in terms of information, postulating different states that give rise to different effects without actually saying what those states are. It is only their position in an information space that counts. If so, then information is a natural candidate to also play a role in a fundamental theory of consciousness. We are led to a conception of the world on which information is truly fundamental, and on which it has two basic aspects, corresponding to the physical and the phenomenal features of the world. (Wikipedia). **Participatory consciousness forms the bastion, pillar, post, stylobate and fulcrum of the redressal measures and corrective steps taken for circumvention of many problems, identified and unidentified, spread widest commonalty. Anne Baring's sententious and pithy article will be good to have a timely outlook.** I would like to open my contribution to this meeting with the words of Ken Wilber that I quoted in the first Millennium Symposium: 'The secret impulse of life is towards greater consciousness. Maybe the evolutionary sequence is from matter to body, to mind, to soul, to spirit - each transcending and including, each with a greater depth and greater

consciousness and wider embrace' Ken Wilber, and with these from Rupert Sheldrake's contribution to the same meeting: "It is possible to make quite a good case that the sun could be conscious and if the sun, then why not the stars? Why not the galaxies? Why should there not be galactic minds? And if they communicate with each other then the distances are far too great for light to be the medium of transmission. We are into inter-galactic telepathy." After many decades of study, I know of only one diagram in all religious traditions that can adequately explain these statements; that can give us both a blue-print of the many levels of consciousness and a route-map to help us reach beyond our present level of understanding. I would like to introduce a diagram of the four worlds of Kabbalah to illustrate the theme that spirit has brought these four worlds into being and that we, living in the fourth and manifest world know nothing of the existence of the other three invisible ones. It is the clearest diagram I know of that can explain to us the multi-levelled structure of consciousness and our place in the great chain of being. We are indeed, as Wilbur suggests, evolving from matter to body, to mind, to soul, to spirit but where have we come from? It is possible that, as participants in the emanation of spirit into manifest life, we have also come from spirit to soul, to mind, to body, to matter. Everything, from the immense galactic energies to the minute forms of matter participates in one life, is generated from one creative spirit. There is nothing outside this life, this spirit. This diagram suggests that truth is not an article of faith but the experience of participating in the undiscovered reality of what has brought us into being. We cannot know truth until we have entered into communion with that reality. Continuing in similar vein **Michel Bitbol dilates upon the dialectic deliberation and polemical argumentation between** hermeneutists and eliminativists. When he formulated the program of Neurophenomenology, Francisco Varela suggested a balanced methodological dissolution of the “hard problem” of consciousness. He shows that his dissolution is a paradigm which imposes itself onto seemingly opposite views, including materialist approaches. I also point out that

Varela's revolutionary epistemological ideas are gaining wider acceptance as a side effect of a recent controversy between hermeneutists and eliminativists. Finally, he emphasizes a structural parallel between the science of consciousness and the distinctive features of quantum mechanics. This parallel, together with the former convergences, point towards the common origin of the main puzzles of both quantum mechanics and the philosophy of mind: neglect of the constitutive blindspot of objective knowledge. Phenomenology and the Cognitive Sciences Phenomenology and the Cognitive Sciences **Phenomenology and the Cognitive Sciences 2002, Volume 1, Issue 2, pp 181-224 Science as if situation mattered Michel Bitbol**

## ABSOLUTE STABILITY

Hence if $X_8 / M_8^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_8$ ) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8 - 8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2 - 2y_8\frac{(E_8 + n_8^2 y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right) - 4y_8(E_8 - n_8)$$

And has the solution

$$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8 - 8)}[y_8(E_8 + n_8^2 y_8) + [y_8^2 (E_8 + n_8^2 y_8)^2 + 4 y_8 (E_8 - 8)(E_8 - n_8)]^{8/2}]$$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2E_8[y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSOILIDATION**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + (M_8 + U_8)\lambda_8 - (M_8 + V_8)\lambda_8 e^{2T_8\lambda_8} + (M_8 U_8 + Y_8)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8\, e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + (M_8 + V_8) Z_8 \sin(Z_8 T_8) - (U_8 M_8 + Y_8) + (M_8(U_8 + V_8) - X_8)$$

$$\mathrm{X}\cos(Z_8 T_8) - 2M_8 V_8 \cos^2(Z_8 T_8) + M_8 V_8 = 0$$

$$(M_8 + U_8) Z_8 - (M_8 + V_8) Z_8 \cos(Z_8 T_8) + (M_8(U_8 + V_8) - X)$$

$$\mathrm{X}\sin((Z_8 T_8) - M_8 V_8 \sin(Z_8 T_8) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 =$

$1 - \cos Z_8\, T_8$and $A_8 = U_8 - V_8$we obtain

$$4M_8^2 V_8^2 W_8^4 + 2M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8(M_8 + V_8) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2[M_8 - V_8)^2(2A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8(M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8(A_8^2 + 4(X_8 + Y_8))]$$

$$+W_8\left[A_8^2+2X_8\right)\left(2A_8M_8-X_8\right)+2A_8Y_8\left(M_8+V_8\right)-M_8A_8^3]+A_8^2\left(X_8M_8-Y_8\right)=0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8>0$$

$$and$$

$$C_8>(X_8-Y_8)/2M_8$$

$recalling\ that\ A_8$= $U_8$and defining $C_8=U_8+V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8(0,U_8V_8=A_8^2(X_8+Y_8)$$

$$Q_8\left(2,U_8V_8\right)=\left(C_8+2M_8^2\right)^2(2C_8M_8+Y_8-X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0⍰$W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$the existence of a D-position at this point is determined by letting cos $(Z_8T_8)=-1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all $T$, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**ABSOLUTE STABILITY IN PREDATOR- PREY MODELS**

**APPLIED TO ACCETUATION ATTRITION MODELS**

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability

criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The Percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending

upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t-T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\big(R(t)\big)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The

regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the

**TIME DELAY**

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a dominant eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

DELAY DUE TO CONGREGATIONAL AND CONFEDERATIONAL DEBIT CREDIT TYEPES

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t-T_8)F_8(R_8(t-T_8) - N_8(t)M_8((R_8(t) \quad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8} + \frac{N_8 dG_8}{dR_8} - \frac{dB_8}{dR_8} > \frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8} + \frac{dM_8}{dR_8}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_8 = \frac{N_8 G_8}{2M_8}\frac{d(8nF_8)}{dR_8} (\geq 0) \quad (8a)$$

$$I_8 = -\frac{N_8 G_8}{2M_8}\frac{d(8nM_8}{dR_8} (\geq 0) \quad (8b)$$

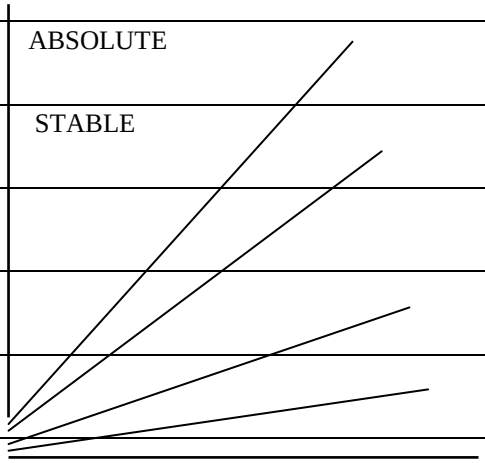
 

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8 > L_8 - I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_8 > L_8 - I_8, \quad \text{when} \quad P_8^2/4 > I \tag{8a}$$

$$P_8 > L_8 - I_8 + \frac{(I_8 - P_8^2}{(L_8 - I)} \quad \text{when} \quad P_8^2/4 < I. \tag{8b}$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit

(like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

See the diagram Nunney gives for predator and prey analysis. We can simulate the same for the financial system in the foregoing too.

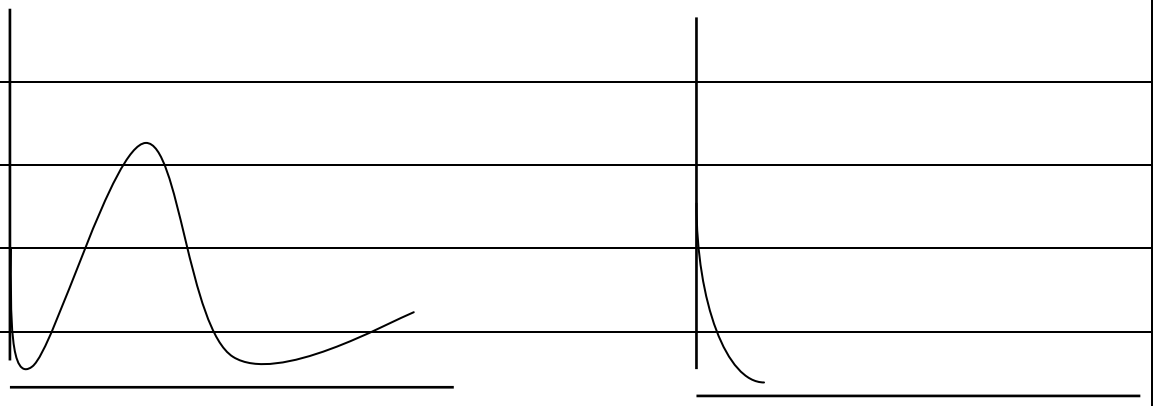
Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the

debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

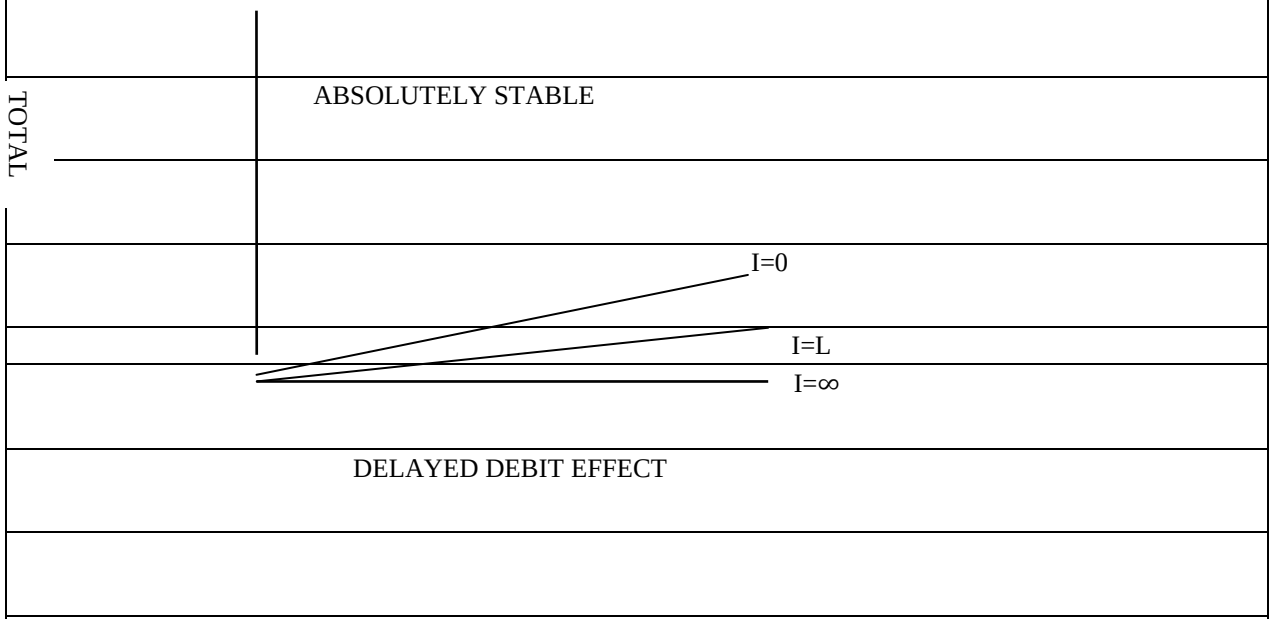
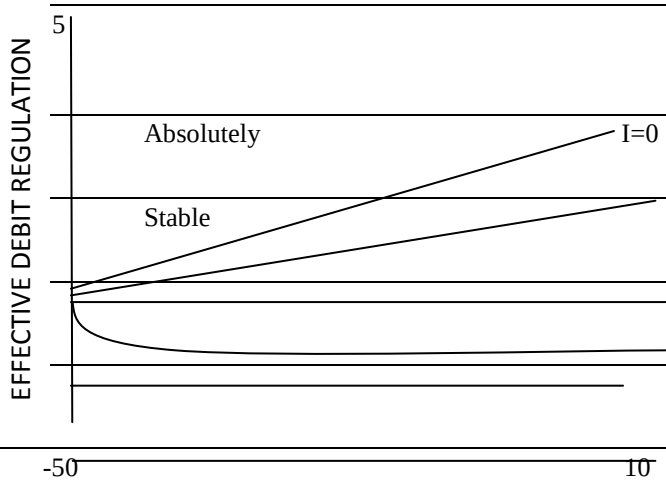



FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison?

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are

defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^{\infty} \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$ then given the constraints of $f(T_8)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8 \lambda_8 e^{-T_8 \lambda_8} + [(U_8 - V_8) M_8 + Y_8] - [((U_8 -$$

$$V_8 M_8 -] e^{-T_8 \lambda_8} = 0$$

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$\boldsymbol{Z}_8^2 + M_8 Z_8 \sin(Z_8 T_8) - (A_8 M_8 + Y_8) + (A_8 M_8 + X_8)\cos(Z_8 T_8) = 0$$

$$(A_8 M_8) Z_8 - M_8 Z_8 \cos(Z_8 T_8) + (A_8 M_8 - X_8)\sin(Z_8 T_8) = 0$$

Where $A_8 = U_8 - V_8$Using (A2.3), substitute for $\boldsymbol{Z}_8^2$ and $Z_8$in (A2.2):

$W_8^2 (X_8^2 - 3 M_8 A_8 X_8 - M_8^2 (X_8 - Y_8) + 2 M_8^3 A_8 + 2 A_8^2 M_8^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_8 = 1 - \cos(Z_8 T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8 / 2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_8 \le 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2\ (\mathrm{X}_8 - Y_8)$$

$$Q_8(2, A_8) = \left(A_8 + 2\,\mathrm{M}_8\right)^2 (2A_8\,\mathrm{M}_8 - (\mathrm{X}_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2\left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8\left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8\left(2X_8 + Y_8\right)\right)$$

$$+2A_8\left(X_8 + Y_8\right)$$

$$=2\,M_8^2 W_8^2 (M_8 + \mathrm{A}_8) + M_8\, A_8^2 W_8 + \left(2A_8 M_8 - X_8 + Y_8\right)(A_8 W_8 + M_8 W_8^2$$

$$+ (2- W_8)((2\,M_8\,X_8 + M_8\,Y_8)\,W_8 + A_8\,(X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $\mathrm{W}_8$touches the $W$ –axis only at its minimum value in which case (A2.9) shows that further increases in *A* will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range $0 \le W_8 \le 2$.Defining

$Q_8\left(W_8, A_8\right) = a_8 W_8^2 + c_8$ *where* $a_8\, b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \square b_8 \square 0$$

And

$$b_8^2 \square 4a_8c_8$$

$$Case(i): Y_8 = 0$$

$Let\ Y_8 = 0$ And try the substitution A= $X_8 = /2M_8 in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8 /$ $2M_8$) is positive for *W<2. Since* $Q_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\ X_8$ ,

$Let\ Y_8 = \ y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 + \ n_8 y_8)$ we have

$$W_8^2\ y_8\left[X_8 n_8 \frac{(1 + n_8 y_8)}{(2 + n_8 y_8)} + M_8^2\left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8 X_8 \left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2}+2y_8\left(n_8-1\right)+\frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)}=0.\ (A2.12)\right.$$

Since $n_8$ and $A_8$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8=(2+n_8 y_8),\ the\ condition(A2.88?)can\ be\ written\ as$

$[\frac{X_8(E_8-8)}{E_8^2}]^2[\frac{X_8^2}{M_8^2}-4y_8 E_8^2\frac{X_8}{M_8^2}+4y_8(n_8-2)E_8^3]$ ?0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8/W_8^2)$ can be expressed as

$2E_8[y_8E_8+(y_8 E_8(8+y_8)(E_8-n_8))^{8/2}]>\frac{X_8}{M_8^2}>2E_8[y_8E_8-y_8E_8(8+y_8)(E_8-n_8))^{8/2}]$

**ABSOLUTE STABILITY**

Hence if $X_8/M_8^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0?$W_8$ ?2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8-8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2-2y_8\frac{(E_8+n_8^2y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right)-4y_8(E_8-n_8)\ ?0$$

And has the solution

$$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8-8)}[y_8(E_8 + n_8^2\, y_8) + [y_8^2\,(E_8 + n_8^2\, y_8)^2 + 4\, y_8 (E_8 - 8)(E_8 - n_8)]^{8/2}]$$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2\, E_8 [y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + (M_8 + U_8)\lambda_8 - (M_8 + V_8)\lambda_8 e^{2T_8\lambda_8} + (M_8 U_8 + Y_8)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8\, e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + (M_8 + V_8) Z_8 \sin(Z_8 T_8) - (U_8 M_8 + Y_8) + (M_8(U_8 + V_8) - X_8)$$

$$\mathrm{X}\cos(Z_8 T_8) - 2M_8 V_8 \cos^2(Z_8 T_8) + M_8 V_8 = 0$$

$$(M_8 + U_8) Z_8 - (M_8 + V_8) Z_8 \cos(Z_8 T_8) + (M_8(U_8 + V_8) - X)$$

$$\mathrm{X}\sin((Z_8 T_8) - M_8 V_8 \sin(Z_8 T_8) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 = 1 - \cos Z_8 T_8$and $A_8 = U_8 - V_8$we obtain

$$4M_8^2 V_8^2 W_8^4 + 2M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8(M_8 + V_8) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2[M_8 - V_8)^2(2A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8(M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8(A_8^2 + 4(X_8 + Y_8))]$$

$$+ W_8[A_8^2 + 2X_8)(2A_8 M_8 - X_8) + 2A_8 Y_8(M_8 + V_8) - M_8 A_8^3] + A_8^2(X_8 M_8 - Y_8) = 0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8 > 0$$

$$and$$

$$C_8 > (X_8 - Y_8)/2M_8$$

$recalling\ that\ A_8 = U_8$ and defining $C_8 = U_8 + V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8(0, U_8 V_8 = A_8^2 (X_8 + Y_8)$$

$$Q_8\left(2, U_8 V_8\right) = \left(C_8 + 2M_8^2\right)^2 (2C_8 M_8 + Y_8 - X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$the existence of a D-position at this point is determined by letting cos $(Z_8 T_8) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all *T*, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**ABSOLUTE STABILITY IN PREDATOR- PREY MODELS**

**APPLIED TO ACCENTUATION ATTRITION MODELS**

We apply model to each and every module and a concatenated model was to be presented which we procrastinate due to spatial restrictions.

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - D(R) - NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

The question is political-discursive. It's what Stengers and Latour are getting at with their concepts of diplomacy and cosmopolitics. They grant, first, that all entities exist and, second, that to say that someone's cherished idol (or whatever disputed entity they hold dear) is non-existent is a **'declaration of war'** – 'this means war,' as Stengers often says. They thus shunt onto-political discourse off of the terrain of knowledge/belief in the sense of existence/non-existence. Their basic claim seems to be that 'respect for otherness,' i.e. political pluralism, can only come from granting the entities that others hold dear an ontology, even if you don't 'believe' in them. You

are thus permitted to say 'I do not follow that god, he has no hold over me' but you are not permitted to say 'your god is an inane, infantile, non-existent fantasy, grow up.' And it's not just a question of politeness (although there's that too). The point is to grant others' idols and deities an existence – one needn't agree over what that existence entails, over what capacities that entity has or what obligations it impresses upon you as someone in its partial presence but to deny it existence entirely is to 'declare war' – to deny the possibility of civil discourse, of pluralistic co-existence. I believe that it was Richard Rorty who once quipped something like claims to reconcile realism and idealism always seem to end with a triumph of idealism. We don't, in fact, get realism through such approaches, but rather just get a pervasive anti-realism. I think this is also the problem with the "non-controversial pluralism" advocated by Stengers and Latour that Phillip defends here. Such pluralism is not realism but is, in fact, a thoroughgoing social constructivism. I think this is the central problem with Latour's argument in Irreductions (these days I regret having ever defended it). In rejecting both Enlightenment critique and what he calls "reduction" he wants to say something like "The Pentacostal really is filled with (e&eb) the Holy Spirit", that for the 19th working scientist heat really is a fluid and phlogiston really is what allows things to burn, and that for the Greek lightning really is an expression of Zeus's anger. Latour tells us that we aren't to reduce or explain away the entities posited by another group's "ontology" but are to develop explanations from within that ontology. It might not sound particularly sexy– and it certainly doesn't tell us what is worth thinking –but I can't help but believe that philosophy is the critical and reflective investigation of basic concepts that guide our investigation of the world about us, how we ought to live our lives, and what form of governance might be best. Compare two figures. A scientist might ask, what causes depression? We can very well imagine a philosopher turning around and asking the scientist, what is causality? The scientist presupposes a concept of causality in her investigations. She uses this concept in her inquiry. Now she might have a sophisticated concept of causality or she might never have thought much about causality at all, using it in the sort of colloquial and unreflective way that Plato decried when, for example, people

like Euthyphro talked about piety. A whole cascade of questions arise when we raise a question like "what is causality?" We can ask whether or not causality exists at all. We can ask how we distinguish between correlation and two events that merely accompany one another from genuine causation. This, for example, was Hume's question. But perhaps most importantly we can ask whether there is only one form of causality or many forms of causality. Is there only one-to-one causation; one cause and one effect? Is there many-to-one causation; or many events conspiring to produce an effect? Is there one-to-many causation; or one event producing a variety of different effects? We can even ask whether causality necessarily moves from past to present or whether there aren't forms of causality that move from future to past! **January 23, 2014 Pluralism and Realism:** multiple-worlds Over at Struggles Forever, Jeremy Trombley has an interesting post up on "the ontological turn" in anthropology or ethnography. I've been meaning to have a discussion with him about this as I think it's an issue many of us are struggling with. For example, the core project of The Democracy of Objects– a project which I think many have missed –is to somehow reconcile some version of social constructivism with a realist ontology capable of making room for ecology (which requires realist and materialist positions as there's a fact of the matter where global warming is concerned) as well as the role played by objective agencies in social assemblages such as technologies, infrastructure, features of geography, local climates, the growth cycles of plants and animals, waste, etc. Maybe we can try to organize some cross-blog event to discuss these issues. I certainly think they're close to the heart of Jeremy, Michael of Archive Fire, Arran James, and a host of others. As an aside, I'm beginning to realize how the different sites of the political I've been outlining– semiopolitics, thermopolitics, oikopolitics (political economy), geopolitics, eropolitics (the politics of sex and desire), biopolitics, and chronopolitics (and I'm sure there are other political sites!) –are drawing me away from traditional Marxism. Assuming that classical Marxism holds that economics or the conditions and relations of production are determinative of all other sites of the political, the various sites of the political that I've been outlining would lead to the conclusion that there is not one determinative

base of the political. This would not require committing Marx to flames, but rather of recognizing the phenomenon of overdetermination, or of a variety of different entangled sites of the political. But I digress. First, I find myself wondering what the ontological turn means in ethnography. Is it 1) the investigation of the different ontologies held/proposed by different cultures? E.g., the Aztecs believed that reality was structured in this way, while the Greeks in that way, and the ancient Chinese this way, etc? Or 2) is it an investigation of how real entities– independent of cultural beliefs –influence cultural formations? Or is it a combination of both? A position that I would favor. January 20, 2014 Model is applied to structured reality

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\left[F\left(R(t)\right) - M\left(R(t)\right)\right]$$

(3)

$$\frac{dR(t)}{dt} = B\left(R(t-T)\right) - D\left(R(t)\right) - N(t)G\left(R(t)\right).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

**The evaporation spectrum of black holes from a local quantum gravity perspective Aurélien Barrau (Submitted on 21 May 2016**)

(5) Authors revisit the hypothesis of (e) a possible line structure in the Hawking evaporation spectrum of black holes, due to (e) non-perturbative quantum gravity effects, even arbitrarily far away from (e) the Planck mass.

(6) Authors show that this naive prediction might in fact hold in (eb)the specific context of loop quantum gravity, with (e&eb) a small departure from (e) the ideal case for (e) some low-spin transitions.

They also show that the effect is (=) neither washed out by the dynamics of (e) the process, nor by (e) existence of a mass spectrum up to (e) a given width, nor by the secondary component induced by (e) the decay of neutral pions emitted during (e&eb) the time-integrated evaporation. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Cite as: arXiv:1605.06613 [gr-qc] (or arXiv:1605.06613v1 [gr-qc] for this version

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t - T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\left(R(t)\right)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

It's very easy to treat this as an abstract, academic debate: "Are you a realist or are you an anti-realist?", as if it were just a matter of what happens to be true. But it seems to me that this debate has, in reality, always been about politics. As I outlined in my last post, we have perpetually seen how appeals to the real and natural have been used in the name of oppressive power, inscribing both the exploitation of nature and the oppression of various people in the very fabric of being itself. Theistic theology and realist ontology have perpetually been used in the name of

what Deleuze called "State Philosophies" or philosophies that ontologize contingent orders of power and privilege (e.g., "the great chain of being" used to justify patriarchy, monarchy, serfdom, poverty, etc, and appeals to nature used to justify poverty and racial inequality (The Bell Curve), patriarchy (evolutionary psychology), heteronormativity, capitalism, etc). Because arrangements of power and inequality are always contingent in the sense that there's no marked difference in the capacities of peoples, power always looks for a transcendent supplement that would provide justification through ontological necessity. Antirealism– from the Greek atomists to present –became the radical and emancipatory gesture because it revealed the lie behind all of these forms of social organization or their inherit contingency or arbitrariness. Realism, by contrast, has all too often functioned as an apologetics for arbitrary power and social organizations. Here it's worth recalling what Foucault said about science: "…Even before we know to what extent something like Marxism or psychoanalysis is analogous to a scientific practice in its day-to-day operations, in its rules of construction, in the concepts it uses, we should be asking the question, asking ourselves about the aspiration to power that is inherent in the claim to being a science. The question or questions that have to be asked are: "What types of knowledge are you trying to disqualify when you say that you are a science? What speaking subject, what discursive subject, what subject of experience and knowledge are you trying to minorize when you say 'I speak this discourse, I am speaking a scientific discourse, and I am a scientist.' What theoretico-political vanguard are you trying to put on the throne in order to detach it from the massive, circulating, and discontinuous forms that knowledge can take?" (Society Must Be Defended, 10). All of these questions hold equally for claims to something being real. What is one trying to minorize when claiming something is real? What becomes privileged? What is excluded? It is these questions that have been at the heart of the realism debates, for as Spencer-Brown taught us, every distinction has a marked and unmarked space, draws attention to something to be included and pushes something into the unconscious or the domain of the invisible, hidden, or veiled. This is above all the case with evocations of the real. However, as I've tried to show antirealism

leads to its own problems. First, so long as we exclude real beings from our ontological inventory, we are unable to fully understand how power functions ("such and such a set of cultural formations have been thoroughly debunked, yet people still live as if they believed them"). Not only do we not fully understand the sources of the problems due to too much focus on the discursive and semiotic, but we deny ourselves valuable sites of political intervention at the level of infrastructure. Second, we are prevented from addressing things such as cultural racism, such as that found in Heidegger's privileging of the West and the Greeks and Germans in particular. We do a good job addressing biological and theological racism, heteronormativity, and sexism by showing how it is a cultural construction, but when faced with racism such as Heidegger's where he argues that there's something "unique" about the Greek event, it's language, and about the German language, or Badiou/Zizek's racism with respect to the "Pauline Event", we really have no response. Here someone like Jared Diamond or Fernand Braud el is needed to explain global-geographical inequalities. Third, the tools of the cultural turn really do not provide us with the means of thinking the ecological as a site of the political. For this reason, I've tried to formulate a third way– which might be called "constructivist realism" or "constructivist naturalism" –that retains the insights of the cultural turn, while also allowing a robust place for the material. I don't claim to be original in this. I think that many such as Manuel DeLanda, Deleuze and Guattari, Stacy Alaimo, Karen Barad, Jane Bennett, etc., are up to something similar. It's a vast project that requires the work of a multitude of voices, especially given the way in which the culturalism pervades humanities. **The Relational and the Non-Relational: Notes towards an Immanent and Pluralist Theory of Meaning Posted by larval subjects under uncategorized k-bigpic I have deleted or reformulated some sentences for which I beg pardon attributable to spatial constraints** What is correlation, what is causation? What is coordination? What is coincidence? These thoughts have bothered humans since long. I have been able to excavate some reliable material from Wikipedia which I would like to share. "All views are partial and contingent – that's the lesson of pluralism, and you say as much yourself when you say that we must be attentive to ways our own

knowledge might contain superstitions. But it's not discourse that determines whether one view works and another doesn't – it's a confrontation with **non-discursive (non-human)** agencies." I wonder if part of the issue here is that we understand pluralism differently. For me this doesn't sound like pluralism at all. Rather, it just sounds like our **epistemological condition**. We [hopefully] come to understand those portions of existence we question or investigate. The reason we come to investigate them is largely contingent. Finally, the accounts of these features of existence we give can be mistaken. All of this is perfectly consistent with a monism. I'm not sure why what he outlines above is the lesson of pluralism. It seems to me that every realist knows this. It seems to me that ontological pluralism is something quite different. Ontological pluralism is the thesis that there are many different worlds inhabited by many different entities. Thus, for example, you would have one world; say that of Lucretius, that's only inhabited by atoms and their combinations. You would have another world; say the world of the Mongolian shaman, that's inhabited by spirits that do all sorts of things. The ontological pluralist is saying that these spirits are, that they exist, that they're real. This is something quite different than merely saying that there are partial and contingent points of view on the world. It's this saying that these entities are rather than that some person or groups of people believe that they are that is the nub of the issue. Now, I think part of the issue here is that there's an ambiguity in the term ontology. Ontology can be one of two things. On the one hand, ontology is a group or persons set of beliefs as to what is. Here it's trivially true that there are pluralities of ontologies and the realist readily recognizes this. This is the whole reason there are debates over ontology. Mongolian shamans have their ontology, Europeans theirs, Christian fundamentalists theirs, materialists theirs, etc. When striving to understand and communicate with others it's vital to understand these ontologies because, as rhetoricians like Burke point out, our **beliefs about what is** are among the things that motivate our action. Despite having never seen bacteria I was my hands because I believe there are bacteria and viruses on door handles and whatnot and don't want to get sick. My belief about a particular thing existing is what motivates my action. And who knows, perhaps this belief is as

superstitious as the belief that the crops failed because God was displeased with my community. On the other hand, ontology is a theory **about what is.** It is making a claim that something exists. This is where the rubber hits the road. The ontological pluralist seems committed to the thesis that every group set of beliefs about what exists is sufficient for granting the existence of those entities. This is what I find objectionable in Latour. Obviously I'm not bothered by Latour's suggestion that we should take into account the role that nonhumans like speed bumps, rivers, microbes, etc., play in the form that social assemblages take. But this is not all that Latour claims. This is precisely where the philosopher might balk. Unlike the ethnographer, the philosopher is not interested in what people believe exists, but rather philosophers– at least of the realist variant –are trying to figure out what is. In other words, the realist philosopher begins with the premise that not all of these beliefs about what is are true. So for the philosopher, recognizing that Mongolian shaman's believe in the existence of shamans would only be the first step. The next step would consist in determining whether there's good reason for thinking such entities really do exist, i.e., whether there's good reason for believing these entities have mind and culture independent reality. Lest readers think that I'm just picking on the supernatural here, we can ask similar questions about strings, subatomic particles, galaxies, etc., etc., etc. I guess the question really comes down to what exactly we mean by ontological pluralism. When we talk about ontological pluralism are we defending the thesis that people have different theories of being? If so, then ontological pluralism is trivially true. If this is what is meant, then I certainly share Jeremy's view that it's valuable to understand the different world's people believe in. Certainly when I was practicing as an analyst I didn't get in ontological debates with my patients and it was necessary to understand their theory of being or their ontology to properly attend to them. Or, when we talk about ontological pluralism are we defending the thesis that all these ontological theories are true and refer to really existing entities? That's quite a different claim and is not one I would defend or endorse. Now someone might object that "in both The Democracy of Objects and Onto-Cartography you defend the thesis that there are multiple worlds." This is true. Because I hold that not everything

is related I'm led to the conclusion that there are diverse worlds. However, I also hold that however many worlds there might be, these worlds are nonetheless composed solely of material entities. Within the framework I propose I wouldn't suggest that there's one world where there are spirits and another world where there are souls and yet another where there are only material entities. My view is that there aren't spirits or souls in any of these worlds. One of the things I keep hearing in these discussions is that somehow the realist adopts a view from nowhere. I honestly don't understand this criticism. Investigation always occurs somewhere and requires all sorts of mediations involving technologies, experiments, etc. It's that labor of gathering evidence, conducting experiments, using technologies to observe the world, etc., that gradually gives us a body of data that allows us to say there's good reason to believe that such and such a thing exists and has these powers. Another charge seems to be that the realist refuses to recognize that their claims about the world are fallible. I find this charge particularly strange because it's precisely because the realist recognizes the difference between our theories of the world or what we say about the world and the world itself that fallibility is built into the core of his position. Realism doesn't mean one holds they have special access to the world, that they know all truths, or that they have the truth in hand, only that there are truths to be known and that we can be mistaken about things. Circling **Squares Philip** has a post responding to my quandaries about how to mesh realism and pluralism. He writes: Ontologically and metaphysically the idea of realist pluralism is no longer an issue. There are (appropriately) numerous variants but the basic idea that reality is itself pluralistic is well established. Captioned expatiation is about Everett worlds and string theory

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from

being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

The von Neumann formulation leads us to conclude that the lower realms of quantum reality form a virtually limitless array of potentials that are inherently incapable of realization without the observation of a knowing entity, which we identify with consciousness (Kafatos and Nadeau, 1990). We thus have a correspondence between quantum theory, the hierarchical structure implied by the Holographic Principle, and the higher orders of experience,

leading to the full expression of the Conscious Universe. David Bohm's theory of the implicate and explicate orders involves a holographic Principle that is fully consistent with the Holographic Principle discussed here (Bohm 1980; 1986; Bohm and Hiley, 1993). According to Bohm, there is an implicate order that represents the universal, holographic subtext of reality, and which unfolds in every moment to produce the explicate order that we all observe. Thomas Germinario (2004) has equated the implicate order with (eb) the unconscious process, and the explicate order with (eb) conscious process. He emphasizes the nature of dreams within (eb) the implicate order, and the importance of the dream work for maintaining a healthy mind through (e&eb) integration of the implicate order and our daily lives. Allan Combs and Mark Holland (1990) connected the implicate order or holomovement (Bohm, 1980) with (e&eb) Carl Jung is theory of synchronicity (Jung and Pauli, 1955), with (e&eb) the implicate order providing (eb) a holographic medium through which apparently disconnected individuals become connected. The principle of synchronicity, the instantaneous connection of people and events beyond (e) the senses, has been equated with the quantum-physical principle of non-locality (Combs and Holland, 1990; Germine, 1991), and has been proposed to be the fundamental mechanism of conscious process (Germine, 1991).Applying the Holographic Principle theory of mind, subjects may be connected synchronously in the manner of two individuals conversing through cellular phones, where the radio signals are transmitted through (e&eb) a distant satellite. The difference here is that the radio signals travel at the speed of light, and thus the there is some miniscule time lapse between the sending and receiving ends. The Holographic Principle implies (eb) a much more distant and instantaneous communication pathway, with (e&eb) the analogue of the satellite being the holographic boundary of the Universe, as well as a much richer communication, potentially involving a transmission of thoughts and feelings. Our ego-consciousness seems to mask (e) the universal relatedness implied by (e) the Holographic Principle, and it is perhaps only through transcendence of (e) the ego-consciousness that the higher orders of experience can become conscious. In the phenomenon of synchronicity, there seems to be a meaningful connection between (e&eb) individuals that

breaks through the barrier of ego-consciousness. Such a connection is reported by many individuals in the course of dreams, when (e) the ego-consciousness has been suspended, at times informing (e&eb) the dreamer of something that has happened in the life of a meaningfully-connected individual. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Impassioned-intense Connectionist neuroscience, in an effort to attain relevance to clinical practice, has recently given us a theory of the relationship of the unconscious to (e&eb) consciousness which is (=) dualistic (Viamontes and Beitman, 2007), an ad hoc theory of conclusions based on (e) assumptions, the most erroneous of which are that representations of (e) information are information, and that such representations elaborate themselves through (e&eb) a process of representations of representations in neuronal circuits (Bechtel and Abrahamsen,1991). As theory, the connectionist models are lacking (e) heuristic value and **are (=) inimical to the progress** of (e) basic and clinical neuroscience. This is not to denigrate or minimize the importance of neuronal connections, but these connections are rather like the wires in (eb) a radio that plays a symphony. The radio, in and of itself, is (=) incapable of playing (e&eb) the symphony. The radio doesn't do not write the symphony, nor does it broadcast the Passionate-raging symphony over the airways. Yet, if we loosen one wire, the symphony is (=) not manifested through the radio. The brain is (=) not subject to (e) the same vulnerability, in that it has (e) parallel and redundant networks of processing, and, it is in this sense that the connectivity of the brain becomes (eb) important. Connectionism, to the extent that it accepts (e) consciousness and the (e&eb) unconscious at all, attributes them to (e) separate groups of "circuits" in the brain, with (e&eb) consciousness-processing circuits having (e) a limited but detailed ability to analyze (e&eb) information, and unconscious-processing processing more information in less detail (Viamontes and Beitman, 2007). There is an element of truth in this view in that information that has importance and meaning resonates within the brain at higher levels of experience and at higher levels of holographic recursion. This resonance involves what we know, what we feel, and what we are capable and willing to consciously realize. The Universe, as William Blake noted,

can be seen in a grain of sand. Indeed, the grain of sand requires (e) the Universe to exist. The human organism requires (e) the Self to exist. The Self is not a simple amalgam of (e) sensory experiences, as suggested in the connectionist model (Blinder, 2007). This amalgam view of the Self can be damaging when applied to (e&eb) psychotherapy, as the connectionists suggest (Blinder, 2007). The Self is ours forever and ever, from eternity to eternity. As Milton suggested, the mind is (=) its own place. Our deepest purpose is for our minds to resonate with (e&eb) the supra-conscious levels of experience, much as it is (=) purpose of the radio to play the symphony. If the symphony is not to our liking, (eb) we can turn the radio off. So, perhaps, the transpersonal consciousness can turn (e&eb) our own individual and collective radios off, or perhaps turn the volume down, if (e) we are not in harmony with the **Universal Mind**. Physics has given us many enigmas. Quantum theory itself is an enigma. But some of the enigmas seem to involve (e&eb) the limitation of (e) our own abilities of discernment. The thought experiment of (e) "Schrödinger's Cat," as it is often interpreted (e.g. Gribbin, 1984), assumes (eb) that the **cat is not conscious, and is (=) incapable of observation in the context of (e) collapse of (e) the wave function.** This is due, in large part, to the false dichotomy of (e) information and (e&eb) experience. Our observation of other animals indicates, to those that are attuned to (e) the animal mind, who can "feel" the animal mind, that, at (eb) the very least, all mammals are (=) conscious. It would seem that the great neuropsychiatrist, Stanley Cobb (1948), may have been correct in his attribution of (e) levels of consciousness to a wide variety of organisms: "...lower animals with (e&eb) no cerebrum appear to be conscious...even plants such as the sunflower that turns towards strong light may have (e) a vague awareness or warmth and comfort. There are many degrees of consciousness and it is my contention that it is integrated at (e&eb) many levels like other important functions of (e) the central nervous system."With respect to the theory of mind, recent mainstream thinking remains (eb) classical and mechanistic. Such views of mind are advocated by our most in prominent neuroscientists (e.g. Changeux, 1985; Gazzaniga, 1985) as neuroscience research attempts to explain all mental function on the basis of (e) **pure mechanism and the localized function of**

**specific areas of the brain.** Francis Crick, the distinguished Nobel Laureate, had a very influential second career as a neuroscientist. He concluded that we are basically soulless creatures, that all mental processes could be reduced to (e&eb) identifiable, neural correlates in (eb) the brain, and that, in fact, we are (=) nothing but a collection of neurons (Crick, 1995). Crick even went so far as to say that what we call the soul is (=) a group of neurons located in the prefrontal lobes of the brain. Another Nobel Laureate, Gerald Edelman (1987), has written a critically acclaimed book on "neural Darwinism," which purports that the brain is "circuitry" is constructed through a developmental process of "survival of the fittest" neurons and neuronal connections. At the same time the total purposelessness and lack of direction in evolution is expounded, again to critical acclaim, by Richard Dawkins (1986). The idea that there essentially no self or soul imbued with agency and self determination was expounded by Gilbert Ryle (1949/2002) with his parody of the "ghost in the machine," and this work continues to be influential among many scholars of the mind. Philosophers such as Patricia Churchland (1986), who coined the term "neurophilosophy," take an "eliminative materialist" view of (e) belief, free will, and consciousness. Other philosophers, such as Daniel Dennett (1991) would make (eb) experience "epiphenomenal," a by-product of (e) brain processes with (e&eb) no effect whatsoever on (eb) the function of mind. However, unlike neuroscience, which is still in its infancy, the philosophy of mind has had a long history, including such notable figures as Plato, Aristotle, Descartes, Kant, Locke, Hume, Kierkegaard, and James, to name just a few. Perhaps, as noted by Ervin Laszlo (1974): "in today's world, most of the traditional functions of (e) cognitive synthesis have atrophied (e) and are ignored and neglected." Hopefully, what we have presented here will be a step in the right direction for what Laszlo calls a "conceptual synthesis" to help fill the need for (e) meaningful engagement is such a world. The fundamental elements of such a synthesis have been described by (e) Laszlo as follows: Conceptual synthesis performs (e&eb) at least five basic functions in the guidance of affairs. They are the mystical, the cosmological, the sociological, the pedagogical or psychological, and the editorial functions. The mystical function inspires (eb)

in man a sense of mystery and profound meaning related to (e&eb) the universe and of himself in it. The cosmological function forms (eb) images of the universe in accord with (e&eb) local knowledge and experience, enabling (eb) men to describe and identify (e&eb) the structure of (e) the universe and the forces of nature. The sociological function validates supports and enforces (eb) social order, representing (eb) it in accord with (e&eb) the nature of the universe, or as the natural or right form of (e) social organization. The pedagogical or psychological function guides (eb) individuals through stages of life, teaching ways of understanding themselves and others and presenting (eb) desirable responses to life is challenges and trials. Finally, the editorial function of (e) conceptual is to define (eb)some aspects of reality as important and credible and hence to be attended to, and other aspects unworthy of serious attention. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Kant sums up this inversion and its spirit early in the Critique of Pure Reason: Up to now it has been assumed that all our cognition must conform to the objects; but all attempts to find something about them a priori through concepts that would extend our cognition have, on this presupposition, come to nothing. Hence let us once try whether we do not get farther with the problems of metaphysics by assuming that the objects must conform to our cognition, which would agree better with the requested possibility of an a priori cognition of them, which is to establish something about objects before they are given to us. This would be just like the first thoughts of Copernicus, who, when he did not make good progress in the explanation of the celestial motions if he assumed that the entire celestial host revolves around the observer, tried to see if he might not have greater success if he made the observer revolve and left the stars at rest. (B xvi) Onticology, like all variations of object-oriented ontology, is realist in its orientation. In defending a realist ontology onticology holds that the vast majorities of objects, actants, beings, or entities are independent of humans and are what they are regardless of whether any humans regard them or register them. In short, onticology rejects any anthropomorphic, idealist, or anti-realist thesis to the effect that to be is to be the correlate of mind, spirit, the body, the human, and language or otherwise. While it is certainly the case that knowledge is necessarily dependent on the object to which it relates, the converse does not hold true. Objects are not dependent on being known, regarded, perceived, or spoken

about. As such, and to put it in Aristotlean terms, knowledge is an accident of objects, not objects an accident of knowledge. As Althusser so nicely puts it, "[n]o doubt there is a relation between thought-about-the-real and thisreal, but it is a relation of knowledge, a relation of adequacy or inadequacy of knowledge, not a real relation, meaning by this a relation inscribed in that real of which the thought is the (adequate or inadequate) knowledge" (Reading Capital, 96). Althusser goes on to remark that "[t]he distinction between a relation of knowledge and a relation of the real is a fundamental one: if we did not respect it we should fall irreversibly into either speculative or empiricist idealism" (ibid.). Onticology categorically endorses Althusser's verdict. It is a fundamental necessity to distinguish between those relations that belong to the object and those that belong to knowledge. Contemporary philosophy continuously confuses these two very different sorts of relations. Naturally the question arises of how it is possible to surmount our relation to the object so as to determine whether objects themselves possess the properties we encounter in relating to objects. In other words, given that we can only ever relate to the object in relating to the object how is it possible to surmount this relation to get at the being of the object itself? Much more will have to be said about this later– and the answers will be surprising with respect to standard prejudices about realism –however, for the moment it can be said that onticology takes its epistemological inspiration from the transcendental realism of Roy Bhaskar. Among other things, Bhaskar sought to provide a transcendental grounding for the sciences. Insofar as onticology defends the thesis that the field of being is much more vast than the field of objects investigated by the natural sciences, it parts way with the thesis that the domain of being is exhausted by the domain of natural objects. However, the general form of Bhaskar's argument holds for our realist purposes. A transcendental argument seeks to elucidate the conditions under which certain acknowledged practices and forms of cognition are possible. Kant, for example, asked what must be the case for mathematical judgments to be possible. How is it both that we are able to extend our knowledge, as if by magic, through mathematical judgments and, more significantly, that these judgments are able to provide genuine knowledge of the world despite the fact that these

forms of reasoning are not based on experience? Part of Kant's argument consisted in claiming that mind imposes the forms of space and time on the data of experience. In other words, space and time are not attributes of being itself but rather of the mind that regards being. Insofar, Kant argues, as mathematics is ultimately a rumination on the nature of space and time taken in their most abstract form and insofar as the mind imposes space and time on the manifold of sensation, it thus follows that a priori judgments about the nature of spatio-temporal relations are possible that anticipate the structure of actual-space times without directly experiencing these space-times. Why? Because any manifold of sensation must necessarily be structured by these forms imposed by intuition. **Onticology–A Manifesto for Object-Oriented Ontology Part I Posted by larval subjects under Object-Oriented Philosophy My Question Is This: You Say, "Granted There Are Deeper And Deeper And More Primitive Forms Of "Knowledge" Or Abstraction Into Mechanism, But I Wouldn't Say That The Atomic Level Brings A New Epistemic World Into View." But Doesn't Any Kind Of Knowledge, However Primitive, Imply Bringing Some Kind Of Epistemic World Into View Or Proto-View?** Joel: David, Yes Exactly. My Question Is, Just How Far Down The Complexity Gradient Can The Meaning Of 'Knowledge,' Or 'Semiotics, Or Specifically Representation (Maya) Be Stretched Without Breaking? I Can Comfortably Stretch It To The Genetic Level, As The Proto-Epistemic, Where Genetic Codes 'Represent' Phenotypic Structures Or Modes, And The Whole Of Evolution Learns From Its Mistakes And Successes As That Knowledge And History Is Encoded Into And As The Recapitulation Of The Embryogenesis Of The Organism. But An Atom, In My Model, While Certainly Reacting To Its Environment From Its Own Attractor, Essence, And 'Interiority' (Prehension) Isn't Making Representations Of It To Do So. But The Key To The Power And 'Essence' Of Representation (And Simultaneously Its Limitation) Is The Capacity Of Abstraction And Choice. Evolution Does This At A Rudimentary Level (And Anthropomorphizing A Bit) With Its Sending Into The World Its 'Random' Variations, Organismic Ideas, In A Sense. Its 'Mind' Is Literally On The Outside As Our Living World And Biosphere. A 'Choice' Is Made In The Biospheric Interiority,

And Knowledge For The Evolutionary/Embryogenetic Trajectory Is Gained By The "Differential Reproductive Success" Of These Individuals Lives Through Time, As The Increasing Intelligence, Representation, And Capacity For Choice Is Directly Injected Into The Flow Through Sexual Selection. But Where Is The Encoded/Abstracted (Enfolded) Representation Of Options And A Choice Among Them Made At The Atomic Level? I Am Working With The Most Cutting-Edge And Coherent Models For The Atom (Including My Own) That I Can Find Among The Heterodoxy, And I See No Evidence Of Representation And Choice At This Level, Or Rather No Real Way To Further Stretch The Meaning Of The Terms To This Level. Sure There Is Uncertainty From Infinite Difference And Immanent Causation (See Bohm On Infinite Causation), But Without Representation There Really Is No Choice (Although Crudely We Could Call A 'Bifurcation' A 'Choice' By The System As A Whole). So I Am Simply Saying That Representation At The Atomic Level Is Entirely Enfolded And These Are Really Just Ontic Level Phenomena. The Atomic Level, In My View, Is Pre- Not Yet Proto-Epistemic, Which I Place at the Genetic Level Where We Find the First Codes And (e&eb) Primitive Representation. This Just Means That There Is A Real Distinction Between Ontic And Epistemic (And A Gradient Between Them) And We Can't Collapse The Two Into Just The Ontic-Epistemic. And Yet, With The Distinction Intact, And The Gradient To Explore, They Remain Nondual, Univocal, ONE Or Ontic. What I Would Say Exists At All Levels (My Own Self-Aware Myth About The Given) Is What I Call The "Symbiogenesis Of Subject And Object". This Is The Spinozan Essence Of Dynamic Stability And Growth, And Leibnizian Prehension. Also Called The "Nucleation Of Observability" In Spinbitz. It Is From This Nucleus Of The Observer (Fuller) In A "Point-Free Geometry" (Whitehead) Of Pointless "Points Of View" As Conditions Of Boundary That Wilber's "All Is Perspective" Finds Its Ontic Grounding. And This Symbiogenesis Is Also A Rudimentary Or Deep-Level Native Or Primitive Intelligence, Evolution At The Involution Of The Subject-Object Interface. It Is An Exploration Of Creative Learning And Direct Immanent Awareness, Unmediated (Enfolded) By Representational Forms (Unfolded). It Just Is, In Waves (Cosmic Vertebrae)

Before It Abstracts, Represents And Incarnates The Flesh Of What Could Be. Matter Is Enfolded Seed Of The Flower Of Its Abstraction. Maya Is This Depth And Recursion Of Brahma Into And Through The Boundary Conditions Of Abstract Relation. And Brahma Is The Span Of The Emergent Rigor Of The Cosmic Ergodic Spine That Opens The Recursions Into The Field Of Infinite Difference. In Immanent-Transcendent Waves, Literally And Empirically A Limit-Cycle In Ergodicity, It Enfolds And Unfolds From Anatom To Anatomy, From Infinite Intelligence Or Omnirelational Awareness Of Self, Through The Representational Depths Of Self As Self-Consciousness. Omni-Evolution, Trans-Dynamics, And The Anatom Posted On February 21, 2013A Discussion With Tom Huston, Part II The Following Is Part 2 Of An Excerpt From A Discussion Initiated By Tom Huston On His Facebook Page, Edited And Embellished For Clarity. Find Part One, Here. In This Part We Dig Into The Nature Of The "Anatom" And The "Ergodic Spine" Of An Electro-Fractal Cosmos, And Into The Trans-Dynamic Integration Between Being And Becoming, And The X-Interface−The Crossroads Of The Ontic-Epistemic (Brahma And Maya) And The Subject-Object Polarities. Univocal_Dynamics_V1 Tom Huston Is A Founding Member Of Integral Institute And A Former Editor Of Enlighten Next Magazine. His Webpage Is Tomhuston.Com. David Marshall Is A Writer And Editor Living In Chicago. No meaning, value, or normativity to be found in nature, that there is nothing **natural beings ought to be,** but that, rather, these judgments arise from us (See last post of larval subjects) which has led some to raise valuable questions about the coherence of these claims. The problem arises when the following three propositions are taken together: 1. there is nothing outside of nature.2. Beings have no intrinsic meaning, purpose, or value (in and of themselves, there's nothing they ought to be). 3. Value judgments about what beings and being ought to be arise from us and beings like us (bonobo apes, dolphins, institutions, birds of paradise, etc). The problem arises between proposition 1 and 3. How can it both be true that there is nothing outside of being and those normative judgments belong to us and other beings capable of making normative judgments, not nature? The problem arises from restricting these judgments to humans and beings capable of

making these judgments. In making such a claim it seems as if we're saying that there's something outside of nature, something that is beyond nature, thereby violating the first thesis and potentially reintroducing the nature/culture distinction. By contrast, if we say that normative judgments are the special domain of those living beings with the proper degree of sentience to make such judgments, then we seem to reintroduce the nature/culture distinction and fall back into the sorts of problems that I outlined in my last post (and that are so nicely critiqued by thinkers such as Latour). Is there a way out of this? I don't know. Great_Chain_of_Being_2_ (lighter) ASIDE: It seems to me that this is really what the debate between realism and anti-realism, realism and socio-linguistic constructivism surrounding the new materialisms and speculative realists has really been all about.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \qquad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9}+\frac{N_9 dG_9}{dR_9}-\frac{dB_9}{dR_9}>\frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9}+\frac{dM_9}{dR_9}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usance bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_9=\frac{N_9 G_9}{2M_9}\frac{d(9nF_9)}{dR_9}(\geq 0) \qquad (9a)$$

$$I_9=-\frac{N_9 G_9}{2M_9}\frac{d(9nM_9}{dR_9}(\geq 0) \qquad (9b)$$

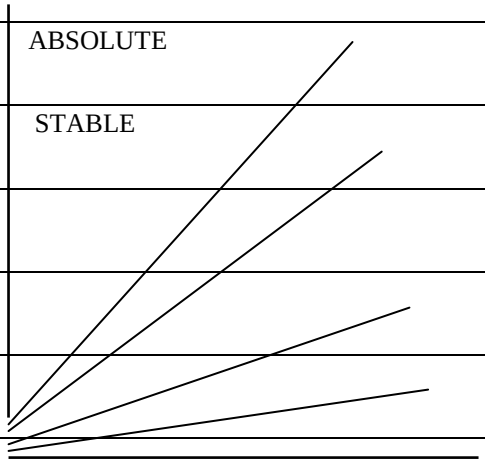
It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9>L_9-I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_9>L_9-I_9$, when $P_9^2/4>I$ (8a)

$P_9>L_9-I_9+\frac{(I_9-P_9^2}{(L_9-I)}$ when $P_9^2/4<I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Reffer couffesy Nunney

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability

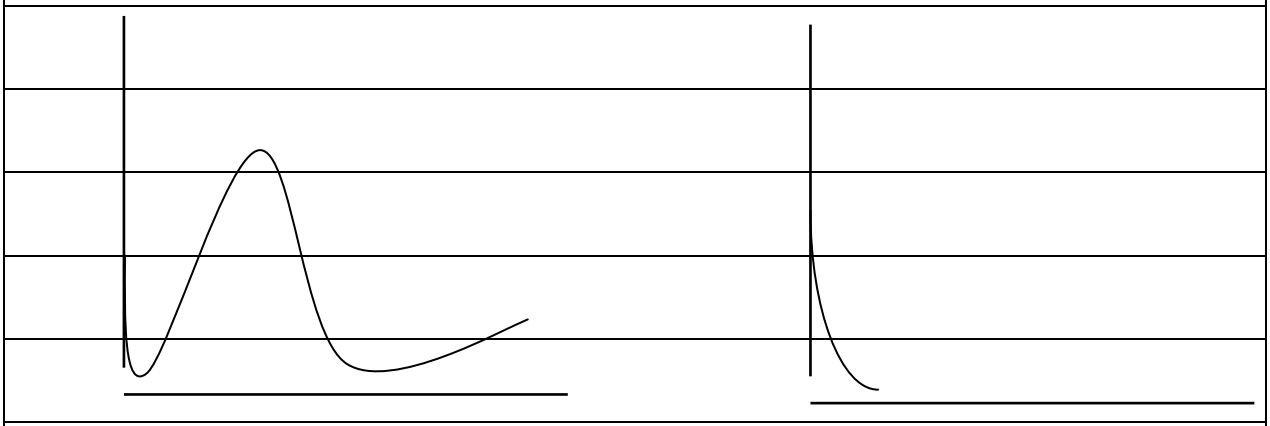
 

is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

Indeed, Nunny's model can be applied to some of the consciousness and philosophical problems. Predator prey analysis has ripped open the veneer of the fundamental concept of universe namely Debit-Credit world.

The noise of ordinary consciousness and (e&eb) beliefs drowns out (e) the signal. Unconsciousness is the background of (e) our ordinary awareness. Our organism is very much at (eb) the center of such effects. The organismic source is (=) our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is a meme-Scape. Pre-conceived concepts vie with structures, concepts with images. Like scientists who ignore assumed truths, we leapfrog over our beliefs and personality deficits, claiming idiosyncratic imagination is literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is both. We inherit the structures that make our experience what it is.

But the structure itself is “empty,” and each human culture “fills” it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us

feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through 're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns, nomenclature, and internal coherence of meaning and representation. In a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist

simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** The lower domain of experience is what we do not directly observe, the quantum realm, while the higher level is what we ordinarily observe, the classical realm. Based purely on experiment, the formulation of quantum theory initially placed the great divide between the quantum and classical domains between the measuring apparatus and the particle (Herbert, 1985). However, in the early 1930s, John von Neumann, in his rigorous mathematical treatise on quantum mechanics, found no support for such a division, compelling him to conclude that the wave for was collapsed by consciousness. Von Neumann's treatise has been called the "quantum bible" and "the most influential book on quantum theory ever written" (Herbert, 1985).

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t-T_9)F_9\big(R_9(t-T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t-T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ and the effective prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting $dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

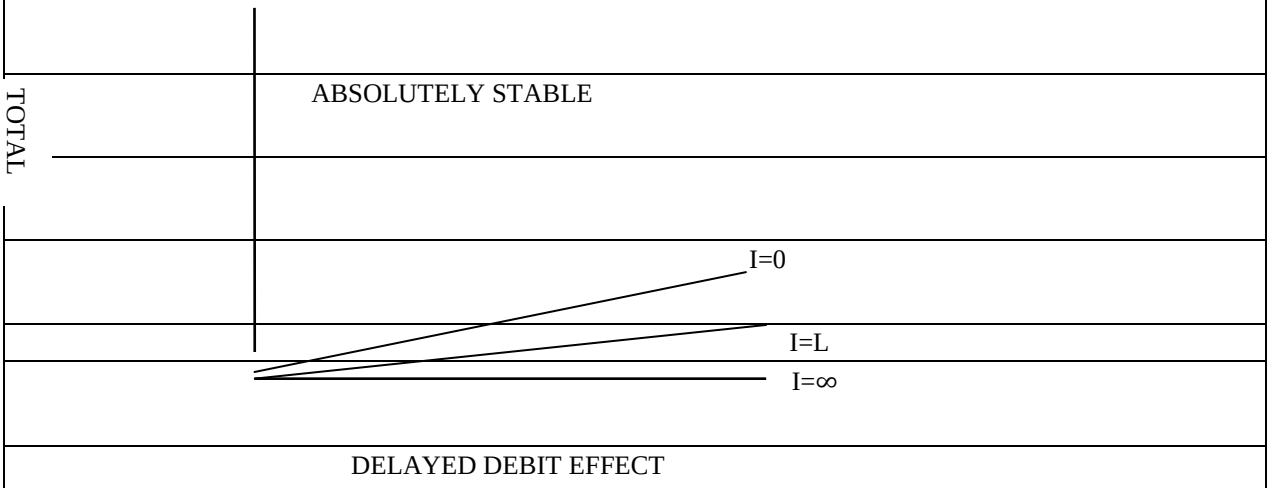



FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.

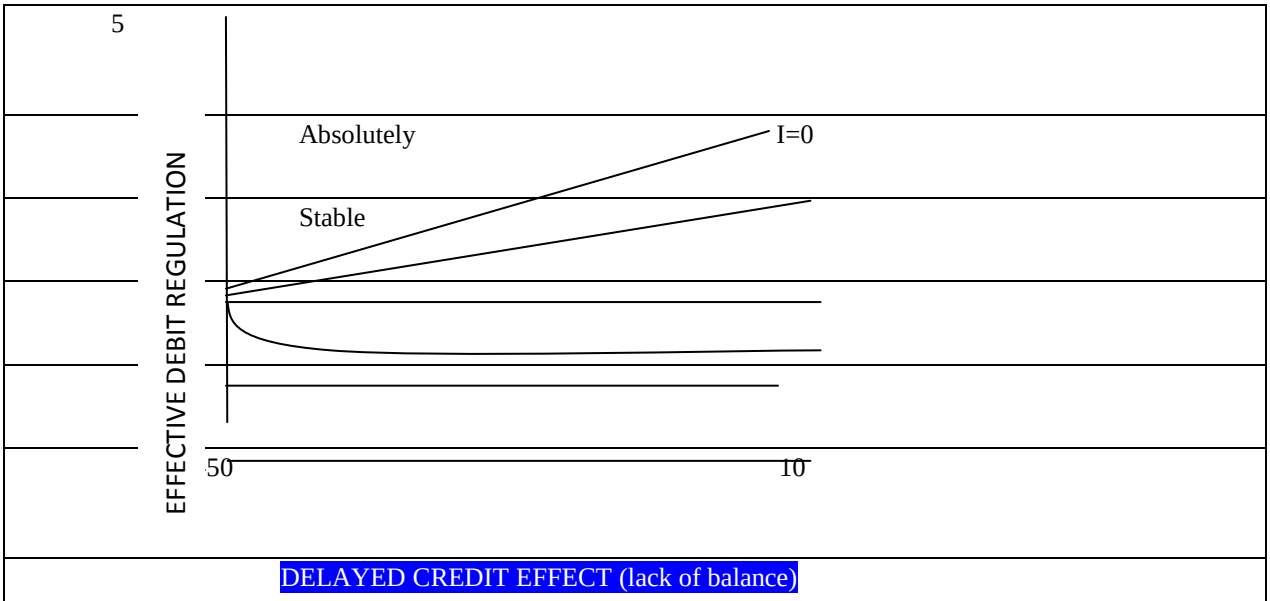



Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the

constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^\infty \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9\left(U_9 - V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9\right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$ then given the constraints of $f(T_9)$, $(A1.5)$ *can never be satisfied* Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDERATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 - V_9M_9 -]e^{-T_9\lambda_9} = 0$$

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9 Z_9 \sin(Z_9 T_9) - (A_9 M_9 + Y_9) + (A_9 M_9 + X_9)\cos(Z_9 T_9) = 0$$

$$(A_9 M_9) Z_9 - M_9 Z_9 \cos(Z_9 T_9) + (A_9 M_9 - X_9)\sin(Z_9 T_9) = 0$$

Where $A_9 = U_9 - V_9$ Using (A2.3), substitute for $Z_9^2$ and $Z_9$ in (A2.2):

$W_9^2(X_9^2 - 3M_9 A_9 X_9 - M_9^2(X_9 - Y_9) + 2M_9^3 A_9 + 2A_9^2 M_9^2]$

$$+W\left[MA^{3}-A^{2}\mathrm{X}+2\mathrm{MA}(2\mathrm{X}+\mathrm{Y})-2X^{2}\right]+A^{2}(\mathrm{X}+\mathrm{Y})=0$$

*where* $W_9$=1-cos ($\mathrm{Z}_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (\mathrm{X}_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ \mathrm{X}_9 - Y_9/2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range 0⍰$W_9$ ⍰2 *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (\mathrm{X}_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\ \mathrm{M}_9)^2\ (2A_9\mathrm{M}_9 - (\mathrm{X}_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9M_9^2 - 3M_9X_9\right) + W_9\left(3M_9A_9^2 - 2A_9X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9 + Y_9\right)$$

$=2M_9^2W_9^2(M_9 + A_9) + M_9A_9^2W_9 + \left(2A_9M_9 - X_9 + Y_9\right)(A_9W_9 + M_9W_9^2$

$+ (2- W_9)((2M_9X_9 + M_9Y_9)W_9 + A_9(X_9 + Y_9))$,

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range 0≤ $W_9$ ≤2.Defining

$Q_9(W_9, A_9) = a_9W_9^2 + c_9$ *where* $a_9 b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

**Harald Atmanspacher, Hans Primas** show that such an epistemic perspective can be consistently complemented by Einstein's ontically oriented position. **Time, Quantum and Information 2003, pp 301-321 Epistemic and Ontic Quantum Realities Harald Atmanspacher, Hans Primas Robert W. Spekkens** present a toy theory that is based on a simple principle: the number of questions about the physical state of a system that are answered must always be equal to the number that are unanswered in a state of maximal knowledge. Many quantum phenomena are found to have analogues within this toy theory. These include the noncommutativity of measurements, interference, the multiplicity of convex decompositions of a mixed state, the impossibility of discriminating nonorthogonal states, the impossibility of a universal state

inverter, the distinction between bipartite and tripartite entanglement, the monogamy of pure entanglement, no cloning, no broadcasting, remote steering, teleportation, entanglement swapping, dense coding, mutually unbiased bases, and many others. The diversity and quality of these analogies is taken as evidence for the view that quantum states are states of **incomplete knowledge** rather than states of reality. A consideration of the phenomena that the toy theory **fails to reproduce, notably, violations of Bell inequalities** and the existence of a Kochen-Specker theorem, provides clues for how to proceed with this research program.DOI: http://dx.doi.org/10.1103/PhysRevA.75.032110 **Evidence for the epistemic view of quantum states A toy theory Phys Rev A 75, 032110 – Published 19 March 2007 Robert W. Spekkens In such specialized and autonomous areas, it is to the eminent and erudite authors to interpret the contradistinctions and contradictions' in the dove tailing explaining. This I feel would be yeoman service in the creation of Nature's general ledger. M. S. Leifer and O. J. E. Maroney** examine the relationship between quantum contextuality (in both the standard Kochen-Specker sense and in the generalized sense proposed by Spekkens) and models of quantum theory in which the quantum state is maximally epistemic. We find that preparation noncontextual models must be maximally epistemic, and these in turn must be Kochen-Specker noncontextual. This implies that the Kochen-Specker theorem is sufficient to establish both the impossibility of maximally epistemic models and the impossibility of preparation noncontextual models. The implication from preparation noncontextual to maximally epistemic then also yields a proof of Bell's theorem from an Einstein-Podolsky-Rosen-like argument. DOI: http://dx.doi.org/10.1103/PhysRevLett.110.120401 Received 25 August 2012Published 20 March 2013 © 2013 American Physical Society **Maximally Epistemic Interpretations of the Quantum State and Contextuality Phys. Rev. Lett 110, 120401 – Published 20 March 2013 M. S. Leifer and O. J. E. Maroney.** Epistemological, onitic and ontological aspects particularly of quantum mechanics have invited

profound thinkers to contribute their mite to the growing literature on the subject, which includes the contributions from the subterranean realm and ceratoid dualism of philosophy like that Chaos theory of Deleuze, Derrida , Kierkegaard and others. **Andrei Khrennikov** shows that the Dirac-von Neumann formalism for quantum mechanics can be obtained as an approximation of classical statistical field theory. This approximation is based on the Taylor expansion (up to terms of the second order) of classical physical variables – maps f: Ω → R, where Ω is the infinite-dimensional Hilbert space. The space of classical statistical states consists of Gaussian measures ρ on Ω having zero mean value and dispersion σ2 (ρ) ≈ h. This viewpoint to the conventional quantum formalism gives the possibility to create generalized quantum formalisms based on expansions of classical physical variables in the Taylor series up to terms of nth order and considering statistical states ρ having dispersion σ2(ρ) = h n (for n = 2 we obtain the conventional quantum formalism). **December 2005, Volume 18, Issue 7, pp 637-650 Date: 28 Nov 2005 Generalizations of Quantum Mechanics Induced by Classical Statistical Field Theory Andrei Khrennikov** It is widely accepted that consciousness or, in other words, mental activity is in some way correlated to the behavior of the brain or, in other words, material brain activity. Since quantum theory is the most fundamental theory of matter that is currently available, it is a legitimate question to ask whether quantum theory can help us to understand consciousness. Several approaches answering this question affirmatively, proposed in recent decades, will be surveyed. It will be pointed out that they make different epistemological assumptions, refer to different neurophysiological levels of description, and adopt quantum theory in different ways. For each of the approaches discussed, these imply both problematic and promising features which will be indicated. **Discrete Dynamics in Nature and Society Volume 2004 (2004), Issue 1, Pages 51-73 http://dx.doi.org/10.1155/S102602260440106X Quantum theory and consciousness: an overview with selected examples Harald Atmanspacher Copyright © 2004 Hindawi Publishing**

**Corporation. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.** Another thought provoking and investigatory article is by Peter G. Lewis et al. Perhaps the quantum state represents information about reality, and not reality directly. Wave function collapse is then possibly no more mysterious than a Bayesian update of a probability distribution given new data. **Peter G. Lewis et al** consider models for quantum systems with measurement outcomes determined by an underlying physical state of the system but where several quantum states are consistent with a single underlying state—i.e., probability distributions for distinct quantum states overlap. Significantly, they demonstrate by example that additional assumptions are always necessary to rule out such a model. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.150404 **Distinct Quantum States Can Be Compatible with a Single State of Reality Phys. Rev. Lett. 109, 150404 – Published 9 October 2012 Peter G. Lewis et al** Communication complexity of a quantum channel is the minimal amount of classical communication required for classically simulating a process of state preparation, transmission through the channel and subsequent measurement. It establishes and perpetuates a limit on the power of quantum communication in terms of classical resources. **Alberto Montina** shows that classical simulations employing a finite amount of communication can be derived from a special class of hidden variable theories where quantum states represent statistical knowledge about the classical state and not an element of reality. This special class has attracted strong interest very recently. The communication cost of each derived simulation is given by the mutual information between the quantum state and the classical state of the parent hidden variable theory. Finally, find that the communication complexity for single qubits is smaller than 1.28 bits. The previous known upper bound was 1.85 bits. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.110501 **Epistemic View of Quantum States and**

**Communication Complexity of Quantum Channels Phys. Rev. Lett 109, 110501 – Published 12 September 2012 Alberto Montina for models see elsewhere .** Stability analysis of what? Quantum state representing reality or our knowledge of reality? Does the quantum state represent reality or our knowledge of reality? In making this distinction precise, **Nicholas Harrigan, Robert W. Spekkens** are led to a novel classification of hidden variable models of quantum theory. They show that representatives of each class can be found among existing constructions for two-dimensional Hilbert spaces. Approach also provides a fruitful new perspective on arguments for the nonlocality and incompleteness of quantum theory. Specifically, **Nicholas Harrigan, Robert W. Spekkens** show that for models wherein the quantum state has the status of something real, the failure of locality can be established through an argument considerably more straightforward than Bell's theorem. The historical significance of this result becomes evident when one recognizes that the same reasoning is present in Einstein's preferred argument for incompleteness, which dates back to 1935. This fact suggests that Einstein was seeking not just any completion of quantum theory, but one wherein quantum states are solely representative of our knowledge. Hypothesis is supported by an analysis of Einstein's attempts to clarify his views on quantum theory and the circumstance of his otherwise puzzling abandonment of an even simpler argument for incompleteness from 1927. **February 2010, Volume 40, Issue 2, pp 125-157 Date: 09 Jan 2010 Einstein, Incompleteness, and the Epistemic View of Quantum States Nicholas Harrigan, Robert W. Spekkens** statistical model of the probabilistic description of physical reality has been studied by many authors. **Andrei Khrennikov posits** a contextual statistical model of the probabilistic description of physical reality. Here contexts (complexes of physical conditions) are considered as basic elements of reality. There is discussed the relation with QM. He proposes a realistic analogue of Bohr's principle of complementarity. In the opposite of the Bohr's principle, **Andrei Khrennikov** principle has no direct relation with mutual exclusivity for

observables. To distinguish his principle from the Bohr's principle and to give better characterization, he changes the terminology and speaks about supplementarity, instead of complementarity. Supplementarity is based on the interference of probabilities. It has quantitative expression trough a coefficient which can be easily calculated from experimental statistical data. There is need not appeal to the Hilbert space formalism and noncommutativity of operators representing observables. Moreover, in our model there exists a pair of supplementary observables which cannot be represented in the complex Hilbert space. There are discussed applications of the principle of supplementarity outside quantum physics. **October 2005, Volume 35, Issue 10, pp 1655-1693 The Principle of Supplementarity: A Contextual Probabilistic Viewpoint to Complementarity, the Interference of Probabilities and Incompatibility of Variables in Quantum Mechanics Andrei Khrennikov Foundations of Physics** Description of the vacuum in Yang-Mills theory and arrive at a physical interpretation of the pseudoparticle solution and the attendant violation of symmetries has been done by **R. Jackiw and C. Rebbi.** It is here vaccuum stability forms the bastion, pillar, post and stylobate of the mechanisms of yang mills Theory and action of a particle. The existence of topologically inequivalent classical gauge fields gives rise to a family of quantum mechanical vacua, parametrized by a CP-nonconserving angle. The requirement of **vacuum stability** against gauge transformations renders the vacua chirally noninvariant. DOI: http://dx.doi.org/10.1103/PhysRevLett.37.172 **Vacuum Periodicity in a Yang-Mills Quantum Theory Phys. Rev. Lett 37, 172 – Published 19 July 1976 R. Jackiw and C. Rebbi.** Three-dimensional Yang-Mills and gravity theories augmented by gauge-invariant mass terms are analyzed. These topologically nontrivial additions profoundly alter the particle content of the models and lead to (eb) quantization of a dimensionless mass-coupling-constant ratio. The vector field excitations become (=) massive, with (e&eb) spin 1 (rather than massless with spin 0), and the mass provides (eb) an infrared cutoff. The gravitation

acquires (e) mass, mediates (e&eb) finite-range interactions, and has (e) spin 2 (rather than being absent altogether); although its mass term is of (e)third derivative order, there are no ghosts or acausalities. DOI: http://dx.doi.org/10.1103/PhysRevLett.48.975 © 1982 The American Physical Society **Three-Dimensional Massive Gauge Theories Phys. Rev. Lett 48, 975 – Published 12 April 1982 S. Deser et al** All scattering amplitudes in (eb) the maximally supersymmetric N=4 super-Yang–Mills theory possess (e) a new, dual superconformal symmetry which extends the previously discovered dual conformal symmetry of MHV amplitudes. To reveal this property authors formulate the scattering amplitudes as (=) functions on (e&eb) the appropriate dual superspace. Rewritten in this form, all tree-level MHV and next-to-MHV amplitudes exhibit manifest dual superconformal symmetry. Authors propose a new, compact and Lorentz covariant formula for (e) the tree-level NMHV amplitudes for (e) arbitrary numbers and types of external particles. The dual superconformal symmetry is broken at loop level by infrared divergences. However, authors provide evidence that the dual conformal anomaly of (e) the MHV and NMHV superamplitudes is (=) the same and, therefore, their ratio is (=) dual conformally invariant. Authors show this explicitly for (e) the six-particle amplitudes at one loop. Authors conjecture that these properties hold for all, MHV and non-MHV, superamplitudes in N=4 SYM both at (eb) weak and at strong coupling. **Nuclear Physics B Volume 828, Issues 1–2, 21 March 2010, Pages 317–374 Dual superconformal symmetry of scattering amplitudes in N=4 super-Yang–Mills theory J.M. Drummonda DOI: 10.1016/j.nuclphysb.2009.11.022.** By comparison with numerical results in (eb) the maximal Abelian projection of lattice Yang–Mills theory, it is argued that the nonperturbative dynamics of (e) Yang Mills theory can be described by (e) a set of fields that take their values in (eb) the coset space SU(2)/U(1). The Yang–Mills connection is parameterized in (eb) a special way to separate (e&eb) the dependence on the coset field. The coset field is then regarded as (=) a collective variable and a method to obtain its effective

action is developed. It is argued that the physical excitations of (e) the effective action may be (=) knot solitons. A procedure to calculate the mass scale of knot solitons is discussed for (e) lattice gauge theories in the maximal Abelian projection. The approach is extended to (e&eb) the SU (N) Yang–Mills theory. A relation between the large N limit and the monopole dominance is pointed out. **Physics Letters B Volume 458, Issues 2–3, 8 July 1999, Pages 322–330 An effective action for monopoles and knot solitons in Yang–Mills theory Sergei V. Shabanov DOI: 10.1016/S0370-2693(99)00612-7.** It is here the quintessential formulation of consciousness occurs. Like a man standing on the threshold of infinity trying to ponder what lies beyond the veil, which separates the seen from unseen a true seeker tries to lift the individual general ledger to the cosmic general ledger. Like in Darwinian evolution, mind also has to evolve from individual consciousness to cosmic consciousness. In Kaishmiri Shaivism one withdraws from the finite to the infinite, but one also goes on an outward journey from the infinite to the finite, because both the finite and the infinite have an intimate connection. The finite is not seen as unreal, but as a symbol of the infinite. There is no real distinction between them. Those two movements constitute Spanda, a key concept in Kaishmiri Shaivism. Spanda is the pulsation of the Absolute in different phases of being. There are no opposites like subject and object, unity or duality, absolute or relative. They are just different phases of the universal vibration of the Absolute. The goal is to realize or be (=) at once infinite and finite. One does not turn away from appearances (like in Advaita Vedanta), but one realizes that the Absolute manifests all things. Spanda, the eternal pulsation of the Absolute, oscillates between a passion to create and dispassion from the created. Through it the Absolute transform itself into all things and then returns back into the emptiness of its undifferentiated nature. In Kaishmiri Shaivism the Absolute is seen as pure consciousness (=being). The Absolute is an eternal all-pervasive principle, the highest reality, the nature of all entities eternally and blissfully at rest within its own nature. The Absolute is the nature of the Self (and

thus of us all). The Absolute is divine; it is Shiva, the Lord of the Universe. It is full of (e) conscious activity through (e&eb)which it generates the universe, and reabsorbs it into (e&eb) itself at the end of each cycle of creation. Thus we speak of monism, as everything resides within (eb) this one absolute consciousness. It sustains (eb) all things, it embraces (e) all things, and it pervades (e&eb) all things. All things are appearances within (eb) the absolute consciousness, but nevertheless real (in contrast to (e) Advaita Vedanta where (e) appearances are seen as unreal or illusionary). All things appear external (out there, outside ourselves), but they do not have (e) a being on their own. They do not exist as (=) separate entities on their own. Everything is contained within (eb) consciousness. What we see as objects are (=) manifestations of consciousness. The event which constitutes (eb) the universe is always internal events happening within (e&eb) consciousness because (e) their essential nature is consciousness itself. If a physical object were totally material, and independent or external to consciousness, it could never be experienced. The universe and (e&eb) consciousness are two aspects of a whole. The universe is (=) an attribute of consciousness which bears consciousness as (=) its substance. Doctrine of Vibration: S.G.Dyczkowski what you "see" is (=) not always what you "get." Many people mistakenly take their own visions literally without (e) "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in (eb) the physical sense but in a psychological sense. The archetypes of (e) the collective unconscious are arrayed behind (e) the scenes of current worldwide conditions, of (e) crisis and confusion. They mirror (e&eb) our own states back at us, whether we perceive them as (=) such or interpret them plausibly or not (Note the essay has been written in model able form and models are given in various papers.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4\,a_9 \square - b_9 \square 0$$

And

$$b_9^{2} \square 4\,a_9 c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9\, in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\, X_9$ ,

$Let\ Y_9 = \ y_9\, X_9$And try the substitution $A_9 = X_9/(2 + A_9\ y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + \ n_9 y_9)$ we have

$$W_9^{2}\, y_9 \left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^{2}\left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9\, X_9 \left[\frac{X_9 (1 + n_9 y_9)}{M_9^{2}(2 + n_9 y_9)^{2}} + 2\, y_9 \left(n_9 - 1\right) + \frac{X_9^{2}(9 + y_9)}{M_9^{2}(2 + n_9 y_9)} = 0.\ (A2.12)\right.$$

Since $n_9$ and $A_9$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9 = (2 + n_9 y_9)$, *the condition*(A2.99?)*can be written as*

$[\frac{X_9(E_9-9)}{E_9^2}]^2[\frac{X_9^2}{M_9^2} - 4\, y_9\ E_9^2 \frac{X_9}{M_9^2} + 4y_9(n_9 - 2)E_9^3]$ .

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9 / W_9^2)$ can be expressed as

$2\,E_9[y_9 E_9 + (y_9\ E_9(9 + y_9)(E_9 - n_9))^{9/2}] > \frac{X_9}{M_9^2} > 2E_9[y_9 E_9 - y_9 E_9(9 + y_9)(E_9 - n_9))^{9/2}]$

Concept Of Being From Shaivite Point Of View: In His Famous Work Entitled On The Conservation Of Force (1847), Hermann Von Helmholtz Had Divided The Totality Of The Energy In The Universe Between Potential Energy And Kinetic Energy And Affirmed The Reciprocal Convertibility Of The Two. In 1852, William Thomson Pointed Out That There Exists A Sub-Ensemble Within Kinetic Energy, Heat, Which, Once It Has Been Generated, Is No Longer Entirely Convertible Into Potential Energy-Or Into Any Other Form Of Kinetic Energy. Considering That The (Partial) Reconversion Of Heat Into Labor Is Possible Only In Situations That Present A Disparity In Temperature, And That Heat Tends To Pass From Warmer To Cooler Bodies By Spreading On An Even Temperature Level Through Space, Thomson Concluded That The Universe Tends Towards A Final State Where Any Energetic Transformations, Every Movement And Every Form Of Life Will Cease: The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO: The Ultimate Nature Of Anything Cannot Change. Change Of Any Sort Is Merely Apparent (Vivarta). The World Of Change And Becoming Is A False Super- Imposition (Adhyaropa, Adhyasa) On The Absolute. It Appears So On The Screen Of Individual, Collective Consciousness, Albeit One Cannot Be So Sure Of Cosmic, For There Are Many Reasons That Cosmic Transactions Also Had To Be Subjected To Change

By ParaBrahaman. In Cosmic Terms, The Mistake (Bhranti) Consists Of The Supposition That The Real Brahman Is The Unreal Universe And The Unreal Universe Is The Real Brahman. In Microcosmic Terms, It Is The Mistake Of Falsely Conceiving The Body, Mind Or Even One's Personality To Be The Self. Note We Have Taken Self As Witness Consciousness. Notwithstanding There Is Nothing Wrong In The Conception Of Self As A Dynamical Presupposition, In That Individual Consciousness And Collective Consciousness And Even Cosmic Consciousness Evolutes In To ParaBrahaman. In The Same Way As The Image Of A Snake Is Falsely Superimposed On A Rope, Similarly The Universe Is Falsely Projected Onto The Real Substratum, The Brahman. Ignorance Is Not Merely A Personal Lack Of Knowledge, But A Cosmic Principle. As Such It Is Called "Maya," The Indefinable Factor (Anirvacaniya) That Brings This Mistake In Identity About. The Reality Status Of This Cosmic Illusion Is Also Indefinable: On The One Hand It Is Not Brahman, The Sole Reality; On The Other Hand It Is Not Absolutely Non-Existent Like A Hare's Horn Or The Son Of A Barren Woman. Note Here The Vedantin Avoids Talking About Zero Which We Have Avidly And Assiduously Stressed About. Nothing Is The Factor That Exists. Also See the Remarks about Satre's Being And Nothingness. Brahman Is The Source Of World Appearances Only In The Sense Of Being Their Unconditioned Ground Or Essential Nature. The Universe Is False Not Because It Has No Nature Of Its Own But Because It Does Have One. The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski. Concept Of Being From Shaivite Point Of View: It Is The 'I Am The Body/Mind' Belief That Gives Rise To The 'I Am Not The World' Belief. These Two Beliefs Are (=) Co-Created."Consciousness Projects The Appearance Of The Mind, Body And World By Taking (Eb) The Shape Of Thinking, Sensing And Perceiving."“Attention Is Consciousness With An Object. When The Object Vanishes, Attention Simply Remains What It Always Is Consciousness."“There Is No Purpose To Meditation. The Purpose Is Already Accomplished."“Everything

That Is Experienced Is Experienced By, Through, In And As Consciousness.”“The Seen Cannot Be Separated From Seeing And Seeing Cannot Be Separated From Consciousness.” “**The Reality Of Any Experience Is Not Hidden In The Appearance; It Is Expressed By The Appearance The Transparency Of Things**' By Rupert Spira Book Review By Dennis Waite The Following Is A Review Of Rupert's Book: 'The Transparency Of Things: Contemplating The Nature Of Experience’. There Is An Essay From The Book Here And The Essay On Cézanne May Be Read At The Awakened Eye Website Nature Rejects The Naiveté That Seeks Absolute Truth. We Are Beginning To Realize, Individually And Culturally, That "Realities" Are All Human Constructions. The Task Becomes One Of "Catching Ourselves In The Act" Of Creating Our Own "Reality" From The Flow Of Events. Human Truth Is Always An Engagement Of Mind With Experience. The Challenge Of The Therapist In These Times Of Chaotic Change Is To Validate The Concept That We Don't Have To Fear The Collapse Of What We Think We Are. Strange Attractors: Transference, Holography, And An **Archetype Burke, J. (2003). Strange Attractors: Transference, Holography, And An Archetype (Doctoral Dissertation, Pacifica Graduate Institute, 2003).** When we get to the topic and understanding of the Transcendental Logic, the brilliant and baffling Transcendental Deduction (TD) Recall the two movements just discussed, the one from experience to its conditions and the one from the forms of valid inference to the concepts that we must use in all judging (the Categories). This duality led Kant to his famous question of right (quid juris) (A84=B116): with what right do we apply the Categories, which are not acquired from experience, to the contents of experience? (A85=B117). Kant's problem here is not as arcane as it might seem. It reflects an important question: How is it that the world as we experience it conforms to our logic? In briefest form, Kant thought that the trick to showing how it is possible for the Categories to apply to experience is to show that it is necessary that they apply (A97).TD has two sides, though Kant never treats them separately John Marlado,

accounts, attributes, and ascribes the philosophy of Nishida KitarÅ is lucid, insightful, and deeply informative despite the highly questionable argument that gives the book its structure. In comparison with other, technically more accurate accounts of Nishida's philosophy, it is broader in scope than most but more single-minded in its approach. The analysis is for the most part limited to four books that have been translated into English, out of approximately 16 volumes of philosophical essays in Japanese, but it aims to capture the contours and most important details of Nishida's entire philosophy. To appreciate the achievement of this account, one would need only to read, in English translation or in the Japanese original, any of the volumes with their streams of repetitive, puzzling expressions and circuitous, often dead- end argumentation. Wilkinson's elucidation of details in the passages he samples makes it easier to understand why Nishida is considered Japan's greatest academic philosopher. A critical look at the frame of the book and some of its own claims reveals where, to my mind, it has gone astray. The argument that frames this book is repeated throughout: Nishida's philosophy is an attempt to articulate, in a characteristically Western philosophical manner, the Zen experience he had as a young man. Nishida first made the attempt by using the conceptual frameworks of Western philosophers, and after they failed him he developed his own, largely by recasting Buddhist insights. In the end, a philosophical articulation of the Zen worldview in Western terms or even in Eastern terms recast by conceptual, Western philosophical methods approved impossible. What Nishida attempted is incommensurable with Western philosophy. While the author first offers these statements as hypotheses (p. 2), one premise of his argument that there is such a thing as "the Zen world-view" or "the Zen experience “remains unquestioned; and another that Nishida's pivotal experience was the Zen experienceâis asserted as "beyond question" (p. 151; see also pp. 28, 48, 10, 114).Wilkinson's presentation of Zen in the first chapter is drawn largely from D. T. Suzuki, with help from Suzuki's portage Abe Masao; some comments on time, birth, and death in; and a few other sources. The

result is a caricature that imagines Zen as an invariable, unitary experience. In fact, Dagan's teachings on practice as the manifestation of enlightenment do not sit well with D. T. Suzuki's emphasis on a satori experience, and Suzuki does not enlighten us about what Zen in its various practices has historically been. The author would have benefited from acquaintance with other presentations by Zen teachers, such as that of ShunryÅ Suzuki'sZen Mind Beginner's Mind, with historical accounts of Zen traditions, and with contrasts between DÅgen's view of time and Sam Â· sÄraand that of other Buddhist philosophers. A little research in contemporary Zen scholarship would have shown how implausible is the talk of "nirvana or awareness of mu" (p. 24), "the basic Zen position [that] never varies" (p. 3), "the Zen vision" (pp. 5, 7), "Zen epistemology" , "Zen experience" (p. 59), "the conception of time involved in Zen" (p. 15), and "the fundamental Zen assertion that there is no absolute time" (p. 16). Even if DÅgen's notion that time and being are nondual (p. 16) epitomized most of Zen, Nishida's early notion of "a transcendent, unchanging reality apart from time" (p. 51) would contrast sharply with it. Equally questionable is Wilkinson's contention that Nishida's "pure experience" is equivalent to Zen satori. The notion of pure experience Nishida developed in his first major work may be pivotal for his entire philosophy, but he never claims or implies that any part of his philosophy presupposes an exceptional experience like satori that grounds his convictions and renders them unverifiable to those who have not had the experience, as Wilkinson intimates (pp. 55, 159, and 161). Nishida presents his notions of pure experience and intellectual intuition, the grasp of its unity, and later... Perennialism and Constructivism is studied by Randolph T Dible : experiencer or the observer and the transcendental subjectivity itself Pure objectivity is essentially a deconstruction, apophatic and negativa The encountering of the experiencer or observer—transcendental subjectivity itself—at the foundation of the world leads inevitably to the recognition of pure objectivity as ultimate reality (which can be taken as its ultimate deconstruction, analogous to the apophatic or via

negativa), from which objects derive their value, weight, significance, meaning or objectivity. In this way, pure objectivity can be seen as the supra-self-evident Axiological Axiom, so to speak, even Unconditional Love, in romantic terms. This axiology (value theory) has a structure inverse to the relationship between transcendental subjectivity as the radical unity of pure self-reference and on the other hand, the world of forms, as mere traces (representations, indications) of the unique, original "first distinction" Spencer-Brown speaks of at the foundation of his calculus. That is, all forms (i.e., distinctions, differences) would reduce to being the first distinction, also known as the marked state, which can be called penultimate reality (pure self-reference or transcendental subjectivity: the Spirit which animates us), except that forms are complimentary to their content, which is their objectivity or value, which would reduce to the unmarked state or ultimate reality. It is the incongruity of form (thoughts; Whitehead's "negative prehensions") and value (feelings; Whitehead's "positive prehensions," or my notion of objectivity, meaning and qualia; in short, the non-formal aspects of experience) that holds forms open and keeps them from absolute reduction. This accounts for the brute, concrete persistence of the "functional illusion"-- to use a term from Dzogchen Buddhism-- of the world. Thus this system has an axiology of metaphysical objectivity grounded on the ideal of pure objectivity as the source of all value, meaning and significance, itself the very fecundity of profundity, which is the motive of drawing the distinction in the first place **(Wikipedia, Kant's writings and Logic Of Sense)** Nano rings were created by accident while intending to make quantum dots. They have interesting optical properties associated with excitons and the Aharonov–Bohm effect. Application of these rings used as light capacitors or a buffer includes photonic computing and communications technology. Analysis and measurement of geometric phases in mesoscopic rings is ongoing. Several experiments, including some reported in 2012, show Aharonov-Bohm oscillations in charge density wave (CDW) current versus magnetic flux, of dominant period h/2e through CDW rings up to 85 **μm** in

circumference above 77 K. This behavior is similar to that of the superconducting quantum interference devices (seeSQUID). **Aharonov–Bohm nano rings (Wikipedia) Stability analysis of nano rings would help progressively accentuate the research about the materials division that uses nano rings.** From the profound contemplation and wisdom of the Buddha and Mahabodhisattvas, we know there is no such reality. Instead, egolessness (non-self) is the only path to understand the reality of the deluded life. All existences are subject to the law of causes and conditions. These include the smallest particles, the relationship between the particles, the planets, and the relationship between them, up to and including the whole universe! From the smallest particles to the biggest matter, there exists no absolute independent identity. Egolessness (non-self) implies the void characteristics of all existence. Egolessness (non-self) signifies (eb) the non-existence of permanent identity for self and existence (dharma). Sunyata stresses (eb) the voidness characteristic of (e) self and existence (dharma). Sunyata and Egolessness possess (e) similar attributes. As we have discussed before, we can observe the profound significance of Sunyata from (e) the perspective of inter-dependent relationships. Considering dharma-nature and (e&eb) the condition of nirvana, all existences are immaterial and of (e) a void-nature. Then we see each existence as (=) independent of each other. But then we cannot find any material that does exist independent of (e) everything else. So Egolessness also implies (eb) void-nature! From the observation of all existences, we can infer (eb) the theory of nirvana and the complete cessation of (e) all phenomena. From the viewpoint of phenomena, all existences are so different from (e) each other, that they may contradict (e&eb) each other. They are (=) so chaotic. In reality, their existence is (=) illusionary and arises from (e) conditional causation. They seem to exist on one hand, and yet do not (e) exist on the other. They seem to be united, but yet they are so different to (e) one another. They seem to exist and yet they do cease! Ultimately everything will return to (e) harmony and complete calmness. This is the nature of all existence. It is the

final resting place for all. If we can understand this reality and remove our illusions, we can find this state of harmony and complete calmness. **Teachings in Chinese Buddhism (6) sunyata (emptiness) in the Mahayana context (Wikipedia) Many models have been given concatenating Buddhist doctrines and Vaccum energy. These models essentially throw light on the essence of self(witness consciousness) and the cosmic general ledger (cosmic consciousness) while the imbalance in the gratification deprivation complex always remain unbalanced when it comes to individual general ledger , advocating the evolutionary theory of individual consciousness to cosmic consciousness. It is to be noted that it is not the question of knowing all the transactions but the wisdom of what lies behind the debit- Credit, Debit-Debit, or Credit-Credit principle that is most important for assimilation.** Look at this simple thesis to prove the point stated in the foregoing. Thesis: 'The World Has A Beginning In Time, And In Space It Is Also Enclosed In Boundaries.' Proof: 'For If One Assumes That The World Has No Beginning In Time, Then Up To Every Given Moment In Time An Eternity Is Elapsed, And Hence An Infinite Series Of States Of Things In The World, Each Following Another, Has Passed Away. But Now The Infinity Of A Series Consists Precisely In The Fact That It Can Never Be Completed Through A Successive Synthesis. Therefore An Infinitely Elapsed World-Series Is Impossible, So A Beginning Of The World Is A Necessary Condition Of Its Existence, Which Was The First Point To Be Proved (Kant, Critique Of Pure Reason, Tr. By R. Guyer And A. W. Wood (Cambridge: Cambridge University Press, 1998): B 454, P. 470). As Quoted In The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO"The Cyclical Hypothesis, So Heavily Criticized By Nietzsche (VP II 325 And 334), Arises In This Way." In Fact, Nietzsche Was Not Criticizing The Cyclical Hypothesis But Only The Particular Form Of That Hypothesis Presented In Vogt's Work. All Of Nietzsche's Texts Without Exception Speak Of The Eternal Return As The Repetitio Of The Same Events Within A Cycle Which Repeats Itself

Eternally. If Deleuze's Interpretation Holds That The Eternal Return Is Not A Circle, Then What Is It? A Wheel Moving Centrifugally, Operating A "Creative Selection," "Nietzsche's Secret Is That The Eternal Return Is Selective" Says Deleuze: The Eternal Return Produces Becoming-Active. It Is Sufficient To Relate The Will To Nothingness To The Eternal Return In Order To Realize That Reactive Forces Do Not Return. However Far They Go, However Deep The Becoming-Reactive Of Forces, Reactive Forces Will Not Return. The Small, Petty, Reactive Man Will Not Return. Affirmation Alone Returns, This That Can Be Affirmed Alone Returns, Joy Alone Returns. Everything That Can Be Denied, Everything That Is Negation, Is Expelled Due To The Very Movement Of The Eternal Return. We Were Entitled To Dread That The Combinations Of Nihilism And Reactivity Would Eternally Return Too. The Eternal Return Must Be Compared To A Wheel; Yet, The Movement Of The Wheel Is Endowed With Centrifugal Powers That Drive Away The Entire Negative. Because Being Imposes Itself On Becoming, It Expels From Itself Everything That Contradicts Affirmation, All Forms Of Nihilism And Reactivity: Bad Conscience, Ressentiment..., We Shall Witness Them Only Once. [...] The Eternal Return Is The Repetition, But The Repetition That Selects, The Repetition That Saves. Here Is The Marvelous Secret Of A Selective And Liberating Repetition. There Is No Need To Remind The Reader That Neither The Image Of A Centrifugal Movement Nor The Concept Of A Negativity-Rejecting Repetition Appears Anywhere In Nietzsche's Writings, And Indeed Deleuze Does Not Refer To Any Text In Support Of This Interpretation. Further, One Could Highlight That Nietzsche Never Formulates The Opposition Between Active And Reactive Forces, Which Constitutes The Broader Framework Of Deleuze's Interpretation. For Some Years, Marco Brusotti Has Called Attention To The Fact That Deleuze Introduced A Dualism That Does Not Exist In Nietzsche's Writings. To Be Sure, The German Philosopher Describes A Certain Number Of "Reactive" Phenomena (For Example, In The Second Essay Of The Genealogy Of Morality, § 11, He Talks About

"Reactive Affects" [Reaktive Affekte], "Reactive Feelings" [Reaktive Gefühlen], Reactive Men [Reaktive Menschen]); But These Are Nonetheless The Result Of Complex Ensembles Of Configurations Of Centers Of Forces That Remain In Themselves Active. Neither The Word Nor The Concept Of "Reactive Forces" Ever Appears In Nietzsche's Philosophy. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO Some Parts Are Deleted Due To Spatial Constraints. In the midst of thoughts about the eternal return we find at least two other thematic axes On the one hand, the view of the **world as a constant flux of forces without any goal, law, or rules of becoming. A chaos sive natura de-divinized and de-anthropomorphized which constitutes the "ontological substratum**" of the whole of Nietzsche's reflexions On the other hand, an ensemble of fragments of an anthropologico-sociological character, designing a path of liberation leading to the creation of superior individuals by way of a profound transformation of their instinctual structure. This transformation must be achieved by a practice of solitude and internal struggle towards the liberation from the ancient representations of the world and from the incorporated herd values. For an analysis of these thematic perspectives, see Paolo D'Iorio, La linea e il circolo. Cosmologia e filosofia dell'eterno ritorno in Nietzsche (Genova: Pantograf, 1995): 233-322. With most of the parameters quantified, with assumptions there must be picturesque view of the self and non self, space and time, science and philosophy with the concatenational consummations. After The Chapter "On Redemption," Where Zarathustra Dares Not Expose His Doctrine, The Eternal Return Begins To Be Enunciated In Part Three Of The Work. In The First Place, It Is The Dwarf Who Formulates It In The Chapter "On The Vision And The Riddle." Facing The "Gate Of The Instant" Which Symbolizes The Two Infinities That Stretch Towards The Past And The Future, The Dwarf Whispers: "All Truth Is Crooked, Time Itself Is A Circle" The Dwarf Represents The Spirit Of Gravity, And He Embodies The Herd Morality, "The Belittling Virtue" Which Is The Title Of Another Chapter From Part III. The Dwarf Can

Endure The Eternal Return Without Great Difficulties Because He Has No Aspirations; Unlike Zarathustra He Does Not Wish To Climb The Mountains That Symbolize Elevation And Solitude. In Two Unpublished Notes, From The Summer And The Fall Of 1883, Nietzsche Writes: The Doctrine Is At First Favored By The Rabble, Before It Gets To The Superior Men. The Doctrine Of Recurrence Will First Smile To The Rabble, Which Is Cold And Without Any Strong Internal Need. It Is The Most Ordinary Of Life Instincts, Which Gives Its Agreement First. Hence, The Content Of The Doctrine Is The Same, But Whereas The Dwarf Can Endure It (Because He Interprets It According To The Pessimistic Tradition For Which "Nothing Is New Under The Sun"), Zarathustra, Who Is The "Advocate Of Life" Regards The Eternal Return As The Strongest Objection To Existence, And As The Rest Of The Dream Suggests, He Does Not Yet Succeed In Accepting It After The Vision At The Gate Of The Instant, The Chapter Is Brought To An End By The Enigma Of The Shepherd. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO .How does the memory affect the GD complex or for that matter the gratification deprivation of an individual or a collective identity? Non-Markovian local-in-time master equations give a relatively simple way to describe the dynamics of open quantum systems with memory effects. Despite their simple form, there are still many misunderstandings related to the physical applicability and interpretation of these equations. Here, we clarify these issues both in the cases of quantum and classical master equations. Further introduction of the concept of a classical non-Markov chain signified through negative jump rates in the chain configuration. **E-M Laine et al 2012 J. Phys. B: At. Mol. Opt. Phys. 45 154004 doi:10.1088/0953-4075/45/15/154004 Local-in-time master equations with memory effects: applicability and interpretation.** Supposing That There Were Indeed An "Energy Of Contraction" Constant In All Centers Of Force Of The Universe, It Remains To Be Explained Where Any Difference Would Ever Originate. It Would Be Necessary For The Whole To Dissolve Into An Infinite Number Of

Perfectly Identical Existential Rings And Spheres, And We Would Therefore Behold Innumerable And Perfectly Identical Worlds COEXISTING [Nietzsche Underlines This Word Twice] Alongside Each Other. Is It Necessary For Me To Admit This? Is It Necessary To Posit An Eternal Coexistence On Top Of The Eternal Succession Of Identical Worlds? The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. Science is usually portrayed as dehumanizing. Brave New World epitomizes this fear. "The more we understand the world, the more it seems completely pointless" (Steven Weinberg). Certainly science can seem chilling when conceived in the abstract as a metaphysical world-picture. We may seem to find ourselves living in a universe with all the human meaning stripped out: Participants in a soulless dance of molecules, or harmonics of pointlessly waggling superstrings and their braneworld cousins. Nature seems loveless and indifferent to our lives. What rights have we to be happy? **(Brave New World: Aldous Huxley).** Religion can afford to be dogmatic but science cannot. Look at this passage which counters both scientists and philosophers alike: However, Hartmann Objects That The Regressive Movement Postulated By Schopenhauer Is Possible Only In Thought: It Remains Nothing More Than An "Ideal Postulate" With No Real Object And Which "Does Not Teach Us Anything About The Real Process Of The World That Unfolds In A Movement Contrary To This Backwards Movement Of Thought" (Hartmann, Philosophie Des Unbewussten, Third Edition (1871): 772). Hartmann Affirms That If Unlike Schopenhauer One Admits The Reality Of Time And Of The World Process, One Must Also Admit That The Process Must Be Limited In The Past And Therefore That There Must Be An Absolute Beginning. In Hartmann's Mind, Failure To Do So Would Result In Positing The Contradictory Concept Of An Accomplished Infinity: "The Infinity That From The Point Of View Of Regressive Thinking, Remains An Ideal Postulate, Which No Reality May Correspond To, Must, For The World, Whose Process Is, On The Contrary, A Progressive Movement, Open Up To A Determinate Result; And Here The Contradiction Comes To Light" (Hartmann

(1871): 772). What Really "Comes To Light" In This Passage Is The Fact That Hartmann Does Not Provide A Demonstration But A Petitio Principii. Indeed, The Concept Of The World Process Analytically Contains The Concept Of A Beginning Of The World. In All Rigors, It Is Therefore Impossible To Demonstrate These Concepts With Reference To Each Other. Secondly, Hartmann's View That One Is Bound To Accept The Reality Of The World Process Even If One Rejects The Ideality Of Schopenhauer's Time Is Mistaken. Hartmann Believes That If Time Is Real There Must Be A World Process With Both An Absolute Beginning And An Absolute End. Without Any Justification, Hartmann Jumps From Schopenhauer's Negated Time To Oriented Time. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. In the last years several theoretical papers discussed if time can be an emergent property deriving from ) quantum correlations. Here, to provide an insight into how this phenomenon can occur, authors present an experiment that illustrates Page and Wootters' mechanism of "static" time, and Gambini et al. subsequent refinements. A static, entangled state between a clock system and the rest of the universe is perceived as evolving by internal observers that test the correlations between the two subsystems. They implement this mechanism using an entangled state of the polarization of two photons, one of which is used as a clock to gauge the evolution of the second: an "internal" observer that becomes correlated with the clock photon sees the other system evolve, while an "external" observer that only observes global properties of the two photons can prove it is static.Phys. Rev. A 89, 052122 (2014)DOI:10.1103/PhysRevA.89.052122 arXiv: 1310.4691 [quant-ph] **Time from quantum entanglement: an experimental illustration Ekaterina Moreva.** Time is an entanglement .When the new ideas of quantum mechanics spread through science like wildfire in the first half of the 20th century, one of the first things physicists did was to apply them to gravity and general relativity. The results were not pretty. It immediately became clear that these two foundations of modern physics were entirely

incompatible. When physicists attempted to meld the approaches, the resulting equations were bedeviled with infinities making it impossible to make sense of the results. Then in the mid-1960s, there was a breakthrough. The physicists John Wheeler and Bryce DeWitt successfully combined the previously incompatible ideas in a key result that has since become known as the Wheeler-DeWitt equation. This is important because it avoids the troublesome infinites—a huge advance. But it didn't take physicists long to realise that while the Wheeler-DeWitt equation solved one significant problem, it introduced another. The new problem was that time played no role in this equation. In effect, it says that nothing ever happens in the universe, a prediction that is clearly at odds with the observational evidence. This conundrum, which physicists call 'the problem of time', has proved to be thorn in flesh of modern physicists, who have tried to ignore it but with little success. Then in 1983, the theorists Don Page and William Wootters came up with a novel solution based on the quantum phenomenon of entanglement. This is the exotic property in which two quantum particles share the same existence, even though they are physically separated. Entanglement is a deep and powerful link and Page and Wootters showed how it can be used (e) to measure time. Their idea was that the way pair of entangled particles evolves is a kind of clock that can be used to measure change. But the results depend on how the observation is made. One way to do this is to compare the change in the entangled particles with an external clock that is entirely independent of the universe. This is equivalent to god-like observer outside the universe measuring the evolution of the particles using an external clock. In this case, Page and Wootters showed that the particles would appear entirely unchanging—that time would not exist in this scenario. But there is another way to do it that gives a different result. This is for an observer inside the universe to compare the evolution of the particles with the rest of the universe. In this case, the internal observer would see a change and this difference in the evolution of entangled particles compared with everything else is an important a measure of time. This is

an elegant and powerful idea. It suggests that time is an emergent phenomenon that comes about because of the nature of entanglement. And it exists only for observers inside the universe. **Time is an entanglement** (Wikipedia). Danish religious philosopher Soren Kierkegaard carried out a systematic critique of the pretensions of reason and an abstract rationalism which he believed that the modern age was nurturing. Condemning reflection as a "danger" that ensnares people in logical delays and machinations, Kierkegaard compared it to a prison. Reflection is for him a form of captivity, a bondage which "can only be broken by [passionate] religious inwardness" (1978: 81). **Reflection seduces individuals into thinking its possibilities are "much more magnificent than a paltry decision"** (1978: 82). It leads them to act **"on principle,"** to dwell on the deliberation of the context of their actions and the calculation of their worth or outcome. Kierkegaard argues that this drives away feeling, inspiration, and spontaneity, all of which are crucial for true inner being and a vital relation to God. For Kierkegaard, as Nietzsche would later agree, genuine inner being (and culture) is characterized by the tautness and tension of the soul which characterizes passionate existence. But the "coiled springs of life relationships ... lose their resilience" in reflection (1978: 78) and "everything becomes meaningless externality, devoid of [internal] character" (1978: 62). Kierkegaard thus contributes to the development of an irrationalist tradition that has echoes in some later postmodern thought. Kierkegaard might well have agreed with his contemporary Fyodor Dostoyevsky, who wrote: "An intelligent [reflective] man cannot seriously become anything ... excessive consciousness is a disease" (1974: 3, 5). In an age overtaken by rules and regulations, genuine action -- which Kierkegaard assumes to be subjective and spontaneous -- is frustrated at every turn. Complaining that we are too "sober and serious" (1978: 71) even at banquets, Kierkegaard bemoans the fact that even suicides are premeditated (1978: 68)! "That a person stands or falls on his actions is becoming obsolete; instead, everybody sits around and does a brilliant job of bungling through with the aid of some reflection

and also by declaring that they all know very well what has to be done" (1978: 73). Thus, it is passion, not reflection, that guarantees "a decent modesty between man and man [and] prevents crude aggressiveness" (1978: 62). "Take away the passion and the propriety also disappears" (1978: 64). The ambiguity in the word "passion" may cause some confusion here. To say that the age and its individuals are "passionless" is not to say there are no emotions whatsoever, but rather that there is no true spiritual inwardness and depth, no intensively motivated action and commitment. It suggests that passion exists only in a simulated, pseudo-form, "the rebirth of passion" through "talkativeness" (1978: 64). "Chattering" for Kierkegaard gets in the way of "essential speaking" and merely "reflects" inconsequential events (1978: 89-99). Hence, in "the present age," emotions -- which in fact are all too pronounced -- have been transformed into negative forces. Anticipating Nietzsche's genealogy of the "slave revolt" in morality, Kierkegaard claims that the "enthusiasm" of the prior age of Revolution, a "positively unifying principle," has become a vicious "envy," a "negatively unifying principle" (1978: 81), a leveling force in its own right insofar as those lacking in talent and resources want to tear down those who have them. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS By Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html.** In November 2011, the United States government argued before the US Supreme Court that it wants to continue utilizing GPS tracking of individuals without first seeking a warrant. In response, Justice Stephen Breyer questioned what this means for a democratic society by referencing Nineteen Eighty-Four. Justice Breyer asked, "If you win this case, then there is nothing to prevent the police or the government from monitoring 24 hours a day the public movement of every citizen of the United States. So if you win, you suddenly produce what sounds like 1984... "[65]. What is to be said in unmistakable terms is that when it comes to national security and innocent citizens there is question of compromise even it be an infringement and encroachment of so called

privacy which has been made a word to subterfuge the activities that are far said to be fair. Quantum theory has provoked intense discussions about its interpretation since its pioneer days. One of the few scientists who have been continuously engaged in this development from both physical and philosophical perspectives is Carl Friedrich von Weizsäcker. The questions he posed were and are inspiring for many, including the authors of this contribution. Weizsäcker developed Bohr's view of quantum theory as a theory of knowledge

## ABSOLUTE STABILITY

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range (0 W 9 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9 - 9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2 - 2y_9\frac{(E_9 + n_9^2 y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right) - 4y_9(E_9 - n_9)$$ 0

And has the solution

$$\frac{X_9}{M_9^2} > \frac{E_9}{(E_9 - 9)}[y_9(E_9 + n_9^2 y_9) + [y_9^2 (E_9 + n_9^2 y_9)^2 + 4y_9(E_9 - 9)(E_9 - n_9)]^{9/2}]$$

Given $E_9 > n_9$ this lower limit to $X_9 / M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9 / M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$ in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9 < n_9$ $(ie., 2A_9 M_9^2 < X_9 - Y_9)$ then (A2.9)and (A2.8)show that $Q_9$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2} > 2E_9[y_9 E_9 - (y_9 E_9 (9 + y_9)(E_9 - n_9)]^{9/2}]$$

And

$$E_9 > n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2 /2M_9 (X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$ and $X_9$ and $Y_9$ by $W_9^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9$ *and* separate (A3.1)into real and imaginary parts

$$Z_9^2 + (M_9 + V_9) Z_9 \sin(Z_9 T_9) - (U_9 M_9 + Y_9) + (M_9(U_9 + V_9) - X_9)$$

$$\mathrm{X}\cos(Z_9 T_9) - 2M_9 V_9 \cos^2(Z_9 T_9) + M_9 V_9 = 0$$

$$(M_9 + U_9) Z_9 - (M_9 + V_9) Z_9 \cos(Z_9 T_9) + (M_9(U_9 + V_9) - X)$$

$$\mathrm{X}\sin((Z_9 T_9) - M_9 V_9 \sin(Z_9 T_9) = 0$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 = 1 - \cos Z_9 T_9$and $A_9 = U_9 - V_9$we obtain

$$4 M_9^2 V_9^2 W_9^4 + 2 M_9 V_9 W_9^3[(M_9 - V_9)^2 + A_9(M_9 + V_9) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2[M_9 - V_9)^2(2 A_9 M_9 + Y_9 - X_9) + (2A_9 M_9 - X_9)^2 - A_9(M_9 - V_9)(2A_9 M_9 - X_9)$$

$$+ M_9 V_9(A_9^2 + 4(X_9 + Y_9))]$$

$$+ W_9[A_9^2 + 2X_9)(2 A_9 M_9 - X_9) + 2A_9 Y_9(M_9 + V_9) - M_9 A_9^3] + A_9^2(X_9 M_9 - Y_9) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9(2, U_9 V_9) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting cos $(Z_9 T_9) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MDDULE NUMBERED SEVEN**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability

criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

**Morphological And Epistemological And Syllogistic Models For Which Nunny's Model Could Be Applied And Has Been Applied:**

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production**

**Charge Asymmetry, $\Delta Y \equiv Yt - Yt^{-}$, Reflects The Asymmetry In $Tt^{-}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles

ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic

quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation

between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to

nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of

Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first

impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of

another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease. Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation

of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces

always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the

thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of

thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity

and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI

Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs” Outcomes—

And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities

Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \quad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

## DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t-T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\big(R(t)\big)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and haimovici model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

**TIME DELAY**

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Limitations of Predator prey analysis in immense. When we realize that world itself is Debit-Credit world and sometime credit-Credit world (See an earlier paper entitled the same nomeneclature.

Opposite limit of absolute stability. The result is a specific The Doctrine of Vibration awareness of an object of the form: 'here this is such.' While individual appearances do not lose their separate identity {svarupabheda) when they rest on a common basis, even so the particular object which appears according to its own characteristics (svalak$ana) is an individual reality in its own right. It is a different kind of appearance characterised by its association with the appearance of the specific location and time in which it is made manifest. The form of our experience is thus 'I now see this here'. But when we perceive each particular constituent appearance separately, each assumes a separate fixed function. Abhinava cites the following colourful example to illustrate how the various combinations of appearances account for the variety of experience: Thus even though the appearance of the beloved may manifest externally, it is as if far away in the absence of another appearance, namely, that of 'embracing'. So when the [appearing of the beloved] is associated with another appearance [namely that of 'far away'] the power (arthakriya) it formerly had of giving pleasure appears as its contrary. The form our experience assumes depends, not only on the nature of the object perceived, but also on personal factors entirely peculiar to ourselves. This theory explains this in two ways. In one sense, the object remains the same, but one or other of its constituent appearances

comes to the fore according to the inclinations of the perceiver. From another point of view, we can say that the perceived object is different for each perceiver according to the difference in the prominent appearance manifest to him. Abhinava, citing as an example a golden jar, illustrates how the same object appears differently to different perceivers according to the use they wish to make of it and to their state of mind: When a person who is depressed and feels that there is nothing [of value for him in the world] sees the jar, he merely perceives the appearance 'exists' [in the form of the awareness that] 'it is\ He is not conscious of any other [of its constituent appearances] at all. An individual who desires to fetch water [perceives] the appearance 'jar'. The man who simply wants something that can be taken somewhere and then brought back [perceives] the appearance 'thing'. The man who desires money [perceives] the appearance 'gold'. The man who desires a pleasing object [perceives] the appearance 'brightness' while he who wants something solid sees the appearance 'hardness'. These 'atomic events' or appearances emerge from the pure subject's consciousness and combine together to form a total event at each moment. Integral Monism of Kashmiri Saivism Daily life (vyavahara) goes on by virtue of this ever renewed flux of appearances. They are connected together and work towards a single unified experience because they appear within the field of consciousness of the universal subject. The aggregate of appearances arises in the [supreme] subject as do [sprouts in] a rice field. Even though each sprout germinates from its own seed, they are perceived as a collective whole. Appearances rest in this way within the universal subject. 'External- ity' is itself another appearance; it arises from a distinction between appearances and the individual subject. So, although all manifestation always occurs within the subject, it appears to be external due to the power of Maya which separates the individual subject from his object. This split must occur for daily life to be possible. Only externally manifest appearances can perform their functions; when they are merged within the subject and at one with him, they cannot do so. Daily life proceeds on the basis of the operation and withdrawal of the conditions necessary for fruitful action to be possible. Appearance in this sense represents the actualisation of a potential hidden in consciousness made possible by virtue of its dynamic, Spanda

nature which is both the flow from inner to outer and back as well as the power that impels it. The emergence from, and submergence into, pure consciousness of each individual appearance is a particular pulsation (visesaspanda) of differentiated awareness. Together these individual pulsations constitute the universal pulse (samanyaspandd) of cosmic creation and destruction. Thus, every single thing in this way forms a part of the radiant vibration {sphuratta, sphurana) of the light of absolute consciousness. Light and Awareness: The Two Aspects of Consciousness Absolute consciousness understood as the unchanging ontological ground of all appearing is termed 'Prakasa'. As the creative awareness of its own Being, the absolute is called ' Vimarsa '. Prakasa and Vimarsa — the Divine Light of consciousness and the reflective awareness this Light has of its own nature — together constitute the all-embracing fullness (purnata) of consciousness. The Recognition (pratyabhijfia) school of Kashmiri Saivism develops this concept of the absolute which finds its fullest expression in Utpaladeva's Stanzas on the Recognition of God. Even though neither of these two key terms appear in the Stanzas on Vibration or the Aphorisms of Siva, they recur frequently in their commentaries. Thus, although the original formulation of the Doctrine of Vibration differs from the theology of Recognition in this respect, it was extended in the course of its development to accommodate this concept of the absolute as well. This was possible, and quite justified, insofar as the absolute understood in Pratyabhijfia terms does not, as we shall see, differ essentially from that of the Spanda school. We can, as Kashmiri Saivites themselves have done, explain one in terms of the other. The Doctrine of Vibration Prakasa: The Light of Consciousness Prakasa is the pure luminosity' (bhdna) or 'self-showing' that constitutes the essence and ultimate identity (atman) of phenomena. That things appear at all is due to the light of consciousness, and their appearing (avabhasana) is itself this Light which bestows on all things their evident, manifest nature. Established in the light of consciousness everything appears there according to its own specific nature (svabhava). Anything that supposedly does not rest in this Light is as unreal as a sky-flower. 3 Thus, according to Rajanaka Rama, unlike the light of the sun, or any other light, this Light not only makes all things apparent, it is also their ultimate source. 4 Full of its

divine vibration the Light makes all things manifest and withdraws them into itself. This supra- temporal activity characterises it most specifically; devoid of it, it would be no better than an inert physical phenomenon. At the same time, this light is the conjunction (slesa) or oneness (aikdtmya) of its countless manifest forms, 6 and the collective whole (Sampina'ana) of all the categories of existence. The universe is nothing but the shining of the Light within itself. It is the radiant vibration (sphuratta) of this Light, the state (avast hand) in which consciousness becomes manifest. Although the Light shines as all things at all times and hence also makes their diversity manifest, 9 penetrating each object individually as well as collectively, it is not totally 'merged' (magna) or identified with the object so as to suffer any division within itself. Our experience of any object is of the form: T see this': it is not itself an object, but the manifest form the object assumes as a luminous principle of experience. The Light is ever revealed and can never be obscured; objectivity can never cast a shadow on the light of consciousness. The Stanzas on Vibration declare: That in which all this creation is established and from whence it arises is nowhere obstructed because it is unconditioned by [its very] nature. This Light is the highest reality (Paramartha). lity with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**DELAY DUE TO CONFEDERATION DEBIT TYPE MATURATION CREDIT TYPE CONSOLIDATION MATURATION**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \quad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9} + \frac{N_9 dG_9}{dR_9} - \frac{dB_9}{dR_9} > \frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9} + \frac{dM_9}{dR_9}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened**

**eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production Charge Asymmetry, $\Delta Y \equiv Yt - Y\bar{t}$, Reflects The Asymmetry In $T\bar{t}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton

particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results

in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely

dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact**

**most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the**

**thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is

nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the

world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease.

Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition

and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the

second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular

chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon

be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining

Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been

Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing

Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

**ABSOLUTE STABILITY**

$$L_9 = \frac{N_9 G_9}{2M_9} \frac{d(9nF_9)}{dR_9} (\geq 0) \quad (9a)$$

$$I_9 = -\frac{N_9 G_9}{2M_9} \frac{d(9nM_9}{dR_9} (\geq 0) \quad (9b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9 > L_9 - I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_9 > L_9 - I_9$, when $P_9^2/4 > I$ (8a)

$P_9 > L_9 - I_9 + \frac{(I_9 - P_9^2}{(L_9 - I)}$ when $P_9^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to

reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

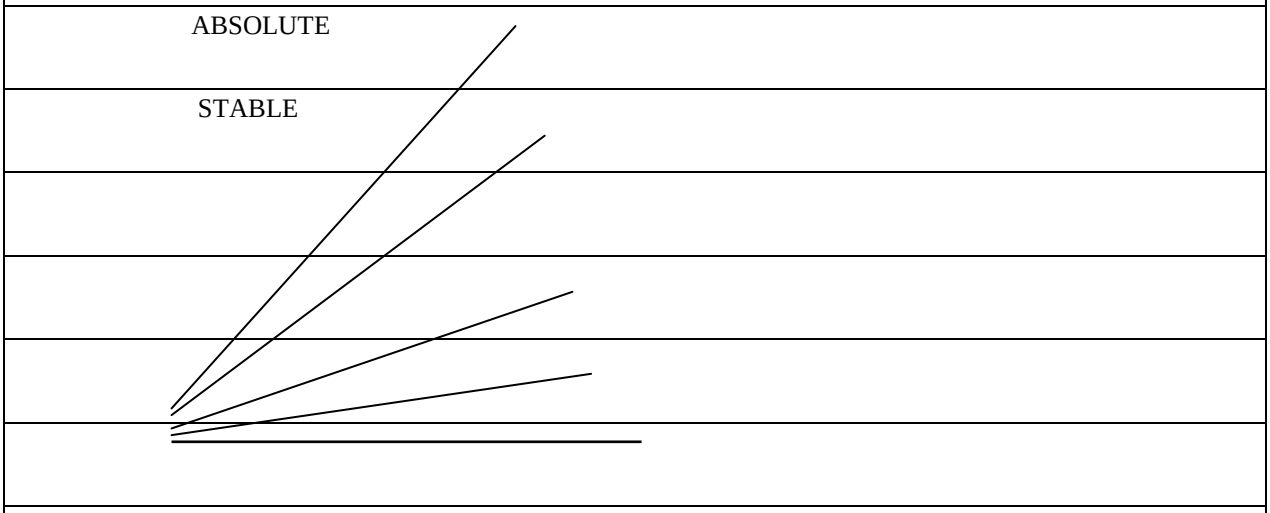



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori

reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

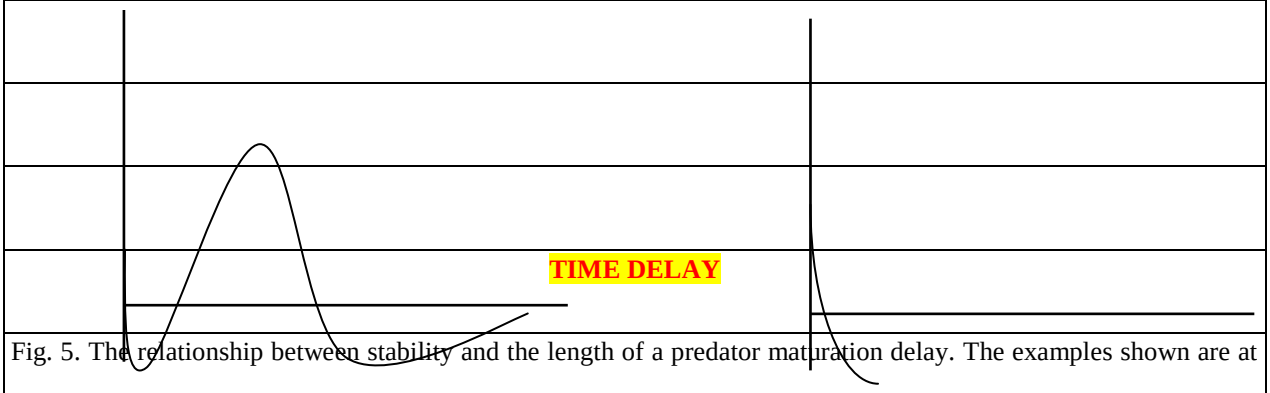



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

## DELAY IN BOTH PREDATOR AND PREY VIS A VIS DEBIT TYPES AND CREDIT TYPES

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t - T_9)F_9\big(R_9(t - T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t - T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) and the effective prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting

$dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

Systems to which model has been applied also include: **Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI:

http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors.

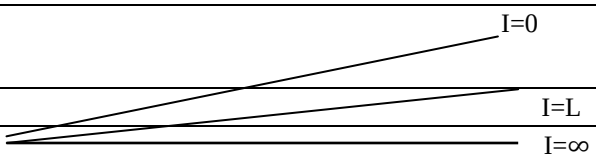
Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, ,**

**Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)**

TOTAL

ABSOLUTELY STABLE

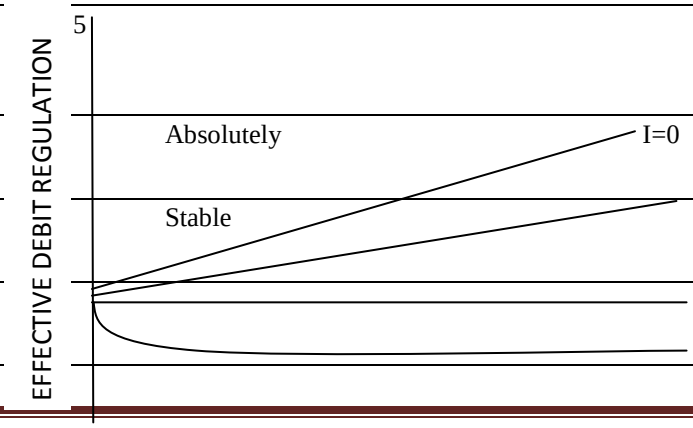



DELAYED DEBIT EFFECT

FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.



-50 10

**DELAYED CREDIT EFFECT (LACK OF BALANCE) CREDIT NON CUMSUMMATION:"**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

The following treatise bears ample testimony and infallible observatory to the fact that how Nunny's model can be applied to various cosmological, general relativity and economic problems.

## ONTOLOGICAL AND EPISTEMOLOGICAL FEATURES OF THE STUDY: OF NUNNY's MODEL

When they are known to be as they are, in the fullest sense of their existence, their phenomenal nature disappears leaving the ground of Being naked and accessible. This approach was validated by a critique of experience. The Vedanta established that space, time and the other primary categories of our daily experience can have no absolute existence. It was therefore necessary to make a distinction between relative truths — that accepted by the precritical common man — and an absolute truth discovered at a higher level of consciousness. The Saiva absolutist 17 rejects any theory that maintains that the universe is less than real. From his point of view a doctrine of two truths, one absolute and the other relative, endangers the very foundation of monism. The Kashmiri Saiva approach is integral: everything is given a place in the economy of the whole. It is equally wrong to say that reality is either one or diverse. Those who do so fail to grasp the true nature of things which is neither as well as both. "We do not" says Abhinavagupta, "base our contention that [reality] is one because of the contradictions inherent in saying that it is dual. It is your approach {paksa) that accepts this [method] While, if [duality and oneness] were in fact [to contradict each other], they would clearly be two [distinct realities]." The Vedantin, who maintains that non-duality is the true nature of the absolute by rejecting duality as only provisionally real, is ultimately landed in a dualism between the real and illusory by the foolishness of his own excessive sophistry (vacafadurvidya). Oneness is better understood as the coextensive unity (ekarasa) of both duality and unity. They are equally expressions of the

absolute. Gopinath Kaviraj says: According to Sarikara, Brahman is truth and Maya is inexplicable (anirvacaniya). Hence the [Advaitin's] endeavour to demonstrate the superiority of Advaita philosophy is turned against his own system. It tarnishes the picture of its philosophical perfection and profundity. He cannot accept Maya to be a reality, therefore his non-dualism is exclusive. The whole system is based on renunciation and elimination and thus is not all-embracing.... By accepting Maya to be Brahman (brahmamayi), eternal (nitya) and real (satyarupa), Brahman and Maya [in the Tantra] become one and coextensive. **The Doctrine of Vibration: An Analysis of the Doctrines and Practices of Kashmir Shaivism Mark S. G. Dyczkowski. Concept of being from Shaivite point of view:** Solutional behaviour and stability analysis of such systems is of paramount importance, notwithstanding the assumptions made which are not farfetched, in that it throws light on the fact whether they could be candidates for yet unsolved problems, providing Autheticative determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication for the wide amplitudinal ramificatory usage of the thesis propounded. The world is now the other person's world, a foreign world that no longer comes from the self, but from the other. The other person is a "threat to the order and arrangement of your whole world...Your world is suddenly haunted by the Other's values, over which you have no control"(Satre: being and Nothingness) Sartre states that many relationships are created by people's attraction not to another person, but rather how that person makes them feel about themselves by how they look at them. This is a state of emotional alienation whereby a person avoids experiencing **their subjectivity by identifying themselves with "the look" of the other**. The consequence is **conflict**. In order to maintain the person's own being, the person must control the other, but must also control the **freedom of the other "as freedom**". These relationships are a profound manifestation of "bad faith" as for-itself is replaced with the other's freedom. The purpose of either participant is not to exist, but to maintain the other participant's looking at them. This system is often mistakenly called "love", but it is, in fact, nothing more than emotional alienation and denial of freedom. In

essence we are always simulating and selling ourselves at the cost of our freedom. Both as a script writer and participant consciousness, human beings experienceâ complete subjective feelings and sometimes absolute subjective feelings, which if remains unchecked might lead to pathological ramifications. (Italics mine) Perspicuous forbearance, Sophisticated seasoning, Partcipational observation, Background actuality, Existential worldliness, Existential strife, Predicational anteriorities, Character consonance, Ontological consonance, Primordial exactitude, Phenomenological correlates, Accolytish representation, Atrophied asseveration. Anamensial alienisms, Anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orinetatilities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state form the bastion, pillar, post, stylobate and sentinel of the classification scheme and doxa. appropriate and diagnostic, differentiating-discriminating, distinctive-distinguishing, emblematic- especial, essentially exclusive, fixed, idiosyncratic-inborn, inbred, indicative and individual, individualistic-inherent, individualizing and ingrained, inherent, innate, local, marked, native, normal, original, particularly- peculiar, personal, private, proper, regular, representative-singular, special, specific, symbolic, symptomatic- unique these form the classification and stratification measure that is suitable notwithstanding the constancy in many a case, as in gravity, which depends on the mass and distance for two given objects**.** Like Kant and Bergson, Deleuze considers traditional notions of space and time as unifying forms imposed by the subject. He therefore concludes that pure difference is non-spatio-temporal; it is an idea, what Deleuze calls "the virtual". (The coinage refers to Proust's definition of what is constant in both the past and the present: "real without being actual, ideal without being abstract.") Proust, Le Temps Retrouvé, ch. III: see the fourth line from the bottom of this page, or, in English translation, the thirteenth paragraph here: "I began to discover the cause by comparing those varying happy impressions which had the common quality of being felt simultaneously at the actual moment and at a distance in time, because of which common quality the noise of the spoon upon the plate, the unevenness of the paving-stones, the taste of the madeleine, imposed the past

upon the present and made me hesitate as to which time I was existing in. Of a truth, the being within me which sensed this impression, sensed what it had in common in former days and now, sensed its extra-temporal character, a being which only appeared when through the medium of the identity of present and past, it found itself in the only setting in which it could exist and enjoy the essence of things, that is, outside Time. [...] **Nothing but a moment of the past Much more perhaps; something which being common to the past and the present, is more essential than both. [...] a marvellous expedient of nature had caused a sensation to flash to me—sound of a spoon and of a hammer, uneven paving-stones**—simultaneously in the past which permitted my imagination to grasp it and in the present in which the shock to my senses caused by the noise had effected a contact between the dreams of the imagination and that of which they are habitually deprived, namely, the idea of existence—and thanks to that stratagem had permitted that being within me to secure, to isolate and to render static for the duration of a lightning flash that which it can never wholly grasp, a fraction of Time in its pure essence. When, with such a shudder of happiness, I heard the sound common, at once, to the spoon touching the plate, to the hammer striking the wheel, to the unevenness of the paving-stones in the courtyard of the Guermantes' mansion and the Baptistry of St. Mark's, it was because that being within me can only be nourished on the essence of things and finds in them alone its subsistence and its delight. It languishes in the observation by the senses of the present sterilised by the intelligence awaiting a future constructed by the will out of fragments of the past and the present from which it removes still more reality, keeping that only which serves the narrow human aim of utilitarian purposes. But let a sound, a scent already heard and breathed in the past be heard and breathed anew, simultaneously in the present and in the past, real without being actual, ideal without being abstract, then instantly the permanent and characteristic essence hidden in things is freed and our true being which has for long seemed dead but was not so in other ways awakes and revives, thanks to this celestial nourishment "While Deleuze's virtual ideas superficially resemble Plato's forms and Kant's ideas of pure reason, they are not originals or models, nor do they transcend possible experience; instead

they are the conditions of actual experience, the internal difference in itself. "The concept they [the conditions] form is identical to its object."[Desert Islands, p. 36.] A Deleuzean idea or concept of difference is therefore not a wraith-like abstraction of an experienced thing, it is a real system of differential relations that creates actual spaces, times, and sensations. [See "The Method of Dramatization" in Desert Islands, and "Actual and Virtual" in Dialogues. In fact it not as if the style and substance of science and philosophers have been entirely different. look at the conformality and consonance in the Gaussian Theory and Deleuze's statements on identity and the loss of it in the following passage…….That is, identity is a continuous process that finds its limit in exhaustion. As Deleuze will write, "…the logical relation of causality is inseparable from a physical process of signalling, without which it would not be translated into action. By 'signal' we mean a system with orders of disparate size, endowed with elements of dissymmetry; by 'sign' we mean what happens within such a system, what flashes across the intervals when a communication takes place between disparates. The sign is indeed an effect, but an effect with two aspects: in one of these it expresses, qua sign, the productive dissymmetry; in the other it tends to cancel it" (dr, 20). An actualized entity just is such a signal-sign system, remaining in perpetual communication with the virtual multiplicity it actualizes, constituting its own elements and identity. Such is Deleuze's account of what Badiou refers to as the operation of the "count-for-one"… an operation that is immanent to being itself and energetic in character. yet how are we to discern the local nature of ontological situations, the manner in which there is no global, overarching global situation in which all local situations are embedded as if parts of a whole? in part the above passage already responds to this question. if series must be brought into relation to communicate in order for causality to occur, and if the resonance of series is always a matter of chance, of a throw of the dice, like Lacan's objet a where no series enjoys primacy of model over copy, then all we have are local situations without a total situation. Unlike Heidegger's analysis of Dasein where Dasein is always being-in-the-world, Deleuze's simulacra or actualizations are always divergent, enjoying only local relations. Heidegger argues that in order to

understand the being of Dasein we must understand the manner in which Dasein is a being-in-the-world. By contrast, Deleuze's concept of multiplicity provides us with the principle of a new structuralism, of a new local ontology, that allows us to understand the immanent organization of a multiplicity without referring it to an embedding global space. as DeLanda so beautifully puts it in relation to the mathematician gauss, "…when gauss began to tap into these differential resources, a curved two-dimension surface was studied using the old Cartesian method: the surface was embedded in three-dimensional space complete with its own fixed set of axes; then, using those axes, coordinates would be assigned to every point of the surface; finally, the geometric links between points determining the form of the surface would be expressed as algebraic relations between the numbers. but gauss realized that the calculus, focusing as it does on infinitesimal points on the surface itself (that is, operating entirely with local information), allowed the study of the surface without any reference to a global embedding space. basically, gauss developed a method to implant the coordinate axes on the surface itself (that is, a method to implant the coordinate axes on the surface itself (that is, a method of 'coordinatizing' the surface) and, once points had been so translated into numbers, to use differential (not algebraic) equations to characterize their relations. as the mathematician and historian Morris Kline observes, by getting rid of the global embedding space and dealing with the surface through its own local properties, 'gauss advanced the totally new concept that a surface is a space in itself'" (intensive science and virtual philosophy, 11-12).according to the old Cartesian method, we can only outline the properties of a space by relativizing it to a global space in terms of which it is then mapped. by contrast, gauss is able to explore a space in terms of its intrinsic metric and organization as a local space, without referencing it to a whole of which it is conceptualized as a part. **The Image of Thought** "Image of thought "permeates both popular and philosophical discourse. According to this image, thinking naturally gravitates towards truth. Thought is divided easily into categories of **truth and error**. The model for thought comes from the educational institution, in which a master sets a problem and the pupil produces a solution which is either true or false. This image of the subject

supposes that there are different faculties, each of which ideally grasps the particular domain of reality to which it is most suited. In philosophy, this conception results in discourses predicated on the argument that "**Everybody knows..."** the truth of some basic idea. Descartes, for example, appeals to the idea that everyone can at least think and therefore exists. Deleuze points out that philosophy of this type attempts to eliminate all objective presuppositions while maintaining subjective ones. Deleuze maintains, with Artaud, that real thinking is one of the most difficult challenges there is. Thinking requires a confrontation with stupidity, the state of being formlessly human without engaging any real problems. One discovers that the real path to truth is through the production of sense: the creation of a texture for thought that relates it to its object. Sense is the membrane that relates thought to its other. Accordingly, learning is not the memorization of facts but the coordination of thought with a reality. "**As a result, 'learning' always takes place in and through the unconscious, thereby establishing the bond of a profound complicity between nature and mind".** Deleuze's alternate image of thought is based on difference, which creates a dynamism that traverses individual faculties and conceptions. This thought is fundamentally energetic and signifying: if it produces propositions, these are wholly secondary to its development. Philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; even tum tan tum;it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the vorticity of being merges with the positive use of the disjunctive synthesis, and

this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoin der for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for every hum of voices or a single silence for all the silences; a single vacuum for all the vacuousness consciousness should not be said without conjuring if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. The attributable to all states of bodies and states of affairs and the expressible of every proposition So univocity of consciousness means the identity of thé noematic attribute and that which is expressed linguistically and sense-fully. Univovity.We mean that it does not allow consciousness to be subsisting in aquasi state and but expresses in all pervading reality. (Reference: LOS: Deleuze) Deleuze sometimes spoke of getting up behind an author and "creating a monster", as if his reading methodology somehow consisted of a distortion of the philosopher's or author's thought. Certainly there is a case to be made for this with regard to the process of deterritorialization, where something is wrested from a territory and deterritorialized upon a new territory, like the animal paw that is deterritorialized from the earth and reterritorialize on the branch. however, it seems to me that the more interesting aspect of Deleuze's approach to other philosophers and art is not so much his "monstrous creations" (i seldom find them particularly monstrous), but rather their Gaussian or Riemannian style, where he explores them in terms of their own internal organization and metric, without reducing them to something alien such as history, society, biology (reductivism), or the signifier. What we find in Deleuze's approach to phenomena is a Gaussian technique. for instance, take Deleuze's books on cinema, his book on Francis bacon, or his study of sacher-masoch. In the first instance, Deleuze's carefully separates cinema from narrative and the signifier, studying it in terms of its specific organization pertaining to the production of images. in the case of bacon, Deleuze doesn't look for an underlying narrative or "meaning", but instead studies the manner in which bacon composes and

organizes his images and lines, both in terms of their production and actuality, so as to liberate a logic of sensation. finally, in approaching sacher-masoch, Deleuze allows Sacher-Masoch's novels to speak for themselves in terms of their desire and relation to pain, stalwartly refusing to reduce masochism to the complement of sadism. in each case, we have a local exploration of a "space" of multiplicities that is extremely precise. Sean Bowden the priority of events: **Deleuze's logic of sense sean Bowden, the priority of events: Deleuze's logic of sense, Edinburgh university press, 2011, 296pp., $40.00 (hbk), isbn 9780748643646 Reviewed by David Scott, Coppin state university and Gaussian spaces and multiplicities a note on Gaussian spaces and multiplicities posted by larval subjects under uncategorized.** What you "see" is not always what you "get." Many people mistakenly take their own visions literally without "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in the physical sense but in a psychological sense. The archetypes of the collective unconscious are arrayed behind the scenes of current worldwide conditions, of crisis and confusion. They mirror our own states back at us, whether we perceive them as such or interpret them plausibly or not. The noise of ordinary consciousness and beliefs drowns out the signal. Unconsciousness is the background of our ordinary awareness. Our organism is very much at the center of such effects. The organismic source is our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is (=) a meme-Scape. Pre-conceived concepts vie with (e&eb) structures, concepts with images. Like scientists who ignore (e) assumed truths, we leapfrog over (e&eb) our beliefs and personality deficits, claiming (eb) idiosyncratic imagination is (=) literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves (e) us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our

experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images. It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is (=) both. We inherit (e&eb) the structures that make our experience (=) what it is. But the structure itself is "empty," and each human culture "fills" it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about (e&eb) our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from

criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns,

nomenclature, and internal coherence of meaning and representation. n a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar some statements in Vedanta, it is not superfluous or redundant. It is evolutionary. Authetication determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication of the evolutionary process towards understanding is ever present in all, its thematic and discursive form. it is the relentless holding on to a theory this is despicable albeit the theory is not. Look at this paragraph. It is the 'I am the body/mind' belief that gives rise to the 'I am not (e) the world' belief. These two beliefs are (=) co-created."Consciousness projects (e&eb) the appearance of the mind, body and world by taking the shape of thinking, sensing and perceiving." "Attention is Consciousness with an object. When the object vanishes, attention simply remains what it always is, Consciousness.""There is (=) no purpose to (e) meditation. The purpose is already accomplished." "Everything that is experienced is experienced by, through, in and as Consciousness." "The seen cannot be separated from seeing and seeing cannot be separated from (e) Consciousness.""The Reality of any experience is not hidden in the appearance; it is expressed by the appearance." "Once we see that everything is Consciousness... Maya still dances, but it is a dance of love not seduction." There are some truly excellent sections

such as a long one taking, as a starting point, the following quotation from the painter Cézanne: "Everything vanishes, falls apart, doesn't it? Nature is always the same but nothing in her that appears to us lasts. Our art must render the thrill of her performance, along with her elements, the appearance of all her changes. It must give us a taste of her Eternity." Rupert begins: "That statement must be one of the clearest and most profound expressions of the nature and purpose of art in our era." And he goes on for some 14 pages to elaborate on this claim, examining the nature of the 'elements' Existence, appearance and Consciousness and their relationship

Some of the following concepts are modeled; some are not. For those which are not modelled kindly see one of the papers in the series. This note has the axiomatic predication to give a holistic view of the study. Here we talk of the characterstics of systems which satisfy the condition of cosmological constant or any other for that matter or variable or axiomatic predication. . There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit,

facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified

from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. Towards the end of circumvention of typing hundreds of superscripts and subscripts which would be a sardine squash and the concomitant operational difficulties, model is presented in piece meal of nine modules. Logarithms are to be taken in respect of those which incorporate more than one variable in bra-ket. Values of log (ab) and log (a+b) are readily available. Anti logarithm shall be taken at the end to predict, project, prognosticate the value of the variable. This is true for tensors, vectors and other variables too. Affirmational assertion, and explanation justification, statement of vindication, annotational commentary, and explicational glossary for each and every system changes and physical interpretation of results is one thing that is to be with earnest endeavour and feverish and febrile expediency. Editions never ever mean the same and identicality of parameters and this has been explicitly and unmistakably stated in the model apriori itself. Akin and analogous, cognate and concurrent, correspondent and congruent, comparable and complementary, synonymous and duple, tantamount and agnate, commensurate and correlative representation is only to highlight the importance and subterfuge the replication of work. It is my fervent solicitation to kindly bear with me for any lapses, notwithstanding the orchestrated efforts for a paper without any minor errors. Postulation predication, conclusive presumption, differential presuppositions, underscored decidedly axiomatic statement of the statement, equation form the bye word or the watch word in the aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement-argument, articulation, assertion, asseveration, choice of words, commentary and communication, declaration and definition, delivery-diction, elucidation-emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation tone and tenor of stratification. Now, how do find the reaction of systems to these singularities. You do the same thing a boss does for you. "Problematize" the events and see how you behave. I will resort to "pressure tactics". "intimidation of deriding report", or "cut in the increment" to make you undergo trials, travails and tribulations. I am happy to see if you improve your work; but may or may not be sad if you succumb to it and hang yourself! We do the same thing with systems. systems show conducive response, felicitous reciprocation or behave erratically with inner roil, eponymous radicality without and with blitzy conviction say like a solipsist nature of bellicose and blustering particles, or for that matter coruscation, trepidiational motion in fluid flows, or seemingly perfidious incendiaries in gormandizing fellow elementary particles, abnormal ebullitions, surcharges calumniations and unwarranted(you think so but the system idoes not!) unrighteous fulminations. Amount of emotion or affect that is appropriate or expected in an given form of interaction. Particular individuals and diffuse obligations (see c and d) are associated with affectivity, whereas contacts with many individuals (universalistic) in a bureaucracy may be devoid of emotion and characterized by affective neutrality. Affective neutrality may refer to self discipline and the deferment of gratification (eg Weber's spirit of capitalism). In contrast, affectivity may be associated with expressing emotions. Adams and Sydie also refer to affective neutrality being associated with ego control (p. 15). (Parsons: Wikipedia). So the point that is made here is "like we problematize the "events" to understand the human behaviour we have to "problematize" the events of systems to understand their behaviour. This statement is made in connection to the fact that there shall be creation or destruction of particles or complete obliteration of the system (blackhole evaporation) or obfuscation of results. Some systems are like "inside traders" they will not put signature atoll! How do you find they did it! Anyway, there are possibilities of a CIA finding out as they recently did! So we can do the same thing witn systems to. This is accentuation, corroboration, fortification, .fomentatory note to explain the various coefficients we have used in the model as also the dissipations called door. In the bank example we have clarified that various systems are individually conservative, and their conservativeness extends holisticallytoo.that one law is universal does not mean

there is complete adjudication of nonexistence of totality or global or holistic figure. Total always exists and "individual" systems always exists, if we donot bring Kant in to picture! For the time being let us not! Equations would become more ensorcelled and frenzied..... philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". In consciousness research, two rival sets of theories can be recognized: (A) Scientific material interpretations of consciousness are based on axioms that view consciousness in the context of highly advanced intentional processing of information in which subject-object relations evolve, and (B) humanistic interpretations of consciousness are based on axioms that view consciousness in the context of, say, "centered pulsations" that enable a conscious agent to act from his or her center of awareness. In this paper I will argue for the selection of axioms that favor humanistic interpretations of consciousness. The Pursuit of Autonomy Interdisciplinary Observations to Human Consciousness: Wautischer, Helmut (1993) the Pursuit of Autonomy. Interdisciplinary Observations to Human Consciousness: Social Neuroscience Bulletin, 6 (4). pp. 52-56.Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality; **CONSCIOUSNESS AND ITS UNIQUENESS June 12, 2012 at 8:40pm (See Deleuze Logic of sense, Wikipedia and Stanford encyclopedia for more details)** Atrocious contrivance and device, gratificational primogeniture, calamitous dodge and expedient, depraved, destructive, disastrous, execrable gambit and gimmick, great solace and succor, iniquitous, injurious, loathsome, low, maleficent, malevolent machination and maneuver, promethaleon of candor, frankness, honesty, honor, ingenuousness, innocence, openness, reality, sincerity, truthfulness, spiteful, stinking, ugly, unpleasant, unpropitious play and ploy, progenitor of jurisprudence and circumspection, racket and ruse, savvy, scam, stratagem, subterfuge, tactic and wile, forthrightness, honesty, truthfulness embodied and personified is the structural predisposition and dispensation of the complex. Nobody need have to take the Brahman-Anti Brahman agency to be a recalcitrant, repugnant, and refractory proposition. God gives and also takes away. That is the point made. ‘

And of course lesser gods, and God doth follow them and they follow God. That is why they are lesser gods. And wield the power on lesser mortals like us. When we mean gratification and deprivation we mean celestial Brahman Anti Brahman and Terrestial Brahman and Anti Brahman agencies producing such happiness or sadness by any means such as (blending- amalgamation, adulterant-adulteration, amalgam-amalgamation, blend- combination,

composite- compound, debasement-denaturant, fusion-hybrid, intermixture-reduction, amalgam- mixture, composite- compound, fusion-mishmash, aggregate- alloy, amalgamation-blend, commixture-composite, composition-compost, conglomerate-fusion, goulash-medley, mishmash-stew, blend-coalition, commixture-compound, federation-heating, integration-junction, melting-merger, synthesis-unification, assortment- combination, adulteration-alloy, assimilation-association, batter-blend, brewed combine, composite compound, concoction confection, conglomeration cross of money, orientation(dis), penury creation, murder, mayhem plunder, pillage, apocalypse, Armageddon in various forms and permutation and combination of Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness: with each action accounted and showing balance in the nature's general ledger as also actionlessness:e). All these are recorded in Nature's general ledger. Neuron DNA encapsulates all the actions passions, actions, interactions, transactions or the lack of it at various levels of individual, collective and cosmic levels. In celestial Brahman Anti Brahman ledger there is cessation of dualities. One more point to be made is if 5=3+2, then we can say 5 gormandizes 3 or 5 gobbles up 2. Quintessentially such a statement arises out of the difference of LHS and RHS term wise. It must be remembered despite all theoretical abstractions and generalizations be it proton mass or Higgs Boson, or topology of a space, it is notwithstanding the fundamentality of nature is just a number. Just everything in Physics and mathematics is just a number. Let us assume I buy rice in Bangalore (District) worth 5k.g. I can always divide 5k.g.s of rice in categories of Bangalore

(South), Bangalore (west), Bangalore (East) and Bangalore (South). On similar basis, there are billions of neutrons, protons and concomitant interactions which would satisfy the given numerical value. Example of gravity above makes things clear. Everything atleast is a function of time and someplace or some other coordinates, subordinates or superordinates. Mass of proton is 2.47 does not mean total mass of all protons is 2.47. It is this argument that is used in the classification scheme. By the very nature, self the unmonitoring but the witness consciousness agency identifies itself with something, be it a profession, an actor on the screen, when it does not have a identity. Such examples are legion, a Charlie Chaplin in "The Kid" or Raj kapoor in "Jagte Raho". These are people who just wanted to live and living without being put under surveillance every moment is the greatest part of their life. This is "Absolute Subjectivity" of Shiva. Such a state is always craving for identification for in pure consciousness all identities are lost. It is this "identitylessness" or having very minimal identity is what pervades most in India and this is due to expansion of individual consciousness. When that happens, be it education, or wealth, the "identity' becomes so well entrenched, they become the exact opposite of "pure consciousness"," impure consciousness" (Shakti). This is a state of "Absolute or relativistic objectivity". There was this case of a lady who went to such an extent of planning, she was sure that whatever she does must work out. When it did not, she lapsed in to "identity crisis". Such thing as this might happen in well developed countries where the competition is extremal and to be achieved at any cost. Realisation that one is not the doer but the witness consciousness is just like watching a movie in which one has acted. It is in this sense "advaita", "non duality' seems to be of cardinal and preeminent importance for the analytical treat ment of "identity' and "identity lessness", both happen in space time. "Multiverses" in this sense is also like a different film of an actor released at different theaters and it is in this sense we state that that every object in space time is a multiverse. You must have totally immersed in doing a "string theory" problem or watching a "film", oblivious of surroundings. Essentially creation (space time) also is what is happening in mind and honed in to you billion times. Identification of anything with a designation or signification is one that leads to

association within spacetime, which would only have pathological consequences and detrimental ramification as stated above or positive implications of "happiness" because of the identity itself. Search or 'show" has to go on! It gives a sense of consummation. a watching pot never boils; donot go to check nature; you are disturbing it; it will never show its true colours bhakti and Virakthi, gratification and deprivation, raga and dwesha is essential to live in space. It is the quantum interest thereof that produces energy to live .gratification and deprivation which form the bastion, pillar, post and stylobate and sentinel of all human actions depends on tams, rajas, and sattva of individual or insentient object. Verily tamas, rajas and satva are the characterstics of space itself, nay the behavioural pattern and attitudinal orientation. There are peaceful places like Vienna or Sringeri and there are violent places like others. Space is warped and woofed by the Akshara, the quantum information of primordial syllable Om. cosmic general ledger is the one vibration of which is brahman; quantum information possessions thereof form the sentinel of all transactions (Gargi to Yajnavalkya: my interpretation)adhocism, nonseqientiality of events, anti sequentiality of events, simultaneity of phenomenon, concurrent happenings , improbability of phenomenon happening itself are characterstics of brahman-anti brahman (Terrestial and celestial) are the cause of non predictability of happenings like in black hole, or in the areas of operation of brahman-anti brahman (Terrestial and celestial). brahman AntiBrahman actions, interactions, decisions, circumspection, adhocism, satisfies sensitive to initial conditions; topologically mixing; dense periodic orbits, the axiomatic predications and postulation alcovishness of chaos theory. Like gravity is a characteristic of space, so is tamas, rajas, sattva, with spaces knotted together by quantum information (for models see below) .accidents are preplanned actions of Terrestial and celestial Brahman AntiBrahman. All actions are perpetrated by them, be it train on fire or death. Belief systems play an important role in the determination of tamas, rajas, and satva state of human beings. There is neither cause nor effect. All actions are performed by Terrestial and celestial Brahman anti Brahman and it is off the record. psychic energy is negative equivalent to –mc^2, so that second law of thermodynamics is satisfied on a holistic basis vis-à-vis material energy

ParaBrahaman is infinite individual, collective and cosmic consciousnesses are infinite (concomitant general ledgers) for they are in the evolution mode with respect to each other in that order. Later emanates from the former as transactions emanate from the conservative state of assets and liabilities. Is not zero true for infinite conservative systems inside space and time? zero is the only consummation, consolidation and conservation. if fifty out of seventy members come out of the room at whatever time you come out, talk about the grand limbo of oblivion and hibernation of lifeless world you are going to and if you are not a celebrity, then it is not coincidence, it is coordination. Nature is playing the same trick on you. That a generalised law is there and you can measure does not make things less good albeit recalcitrant, refractory and rogue particles occur time and again, by and large. Entering in to all living beings through his constant part on the new moon night, he is born therefrom next morning. Absolute ParaBrahaman (zero) projects himself again and again on the screen of individual consciousness which appears of depairing despondent world. Some where the balancing act is done on karma. cosmic general ledger reflects upon itself to keep all transactional ties conservative (Brihadarnkayaka Upanishad: 1.5 14) for models see next paper through its divine power ( i interpret man creates god; it is attribution and ascription of various Shaktis would the formation of god takes place like space and time) the self assumes (eb) the form of a deity man can contemplate and venerate, even though siva, the pure subject, can never in fact be an object of meditation (adhyeya). Until we realise our true identity with Sankara, he is worshipped and conceived to be a reality alien to ourselves. While we are in the realm of creation, he too is a creation or mode or appearing of the absolute, manifest to us in meditation, through his freedom as an eternal, omniscient being. There is no gulf between the created and the uncreated creator: nothing in reality, although an object of knowledge, ceases to be Siva: this is the reason why meditation [on this or that aspect] of reality bestows its fruit. The world of the senses and mind appears to the well awakened (suprabuddha) as a theophany, an eternal revealing of god in his creation. The doctrine of vibration declares that "there is no state in word, meaning or thought, either at the beginning, middle or end, which is not Siva." 59 to utter any word is, in

reality, to intone a sacred formula.every act is a part of Siva's eternal cosmic liturgy, every movement of the body a ritual gesture (mudras), and every thought, god's thought. By what path are you not attainable? What words do not speak of you? in which meditation are you not an object of contemplation? What indeed are you not, o lord? 61 spandashakti, which accounts for the appearing of all things, is also the means by which deity in its many varied forms appears to man. Ksemaraja concludes: the ultimate object of worship of any theistic school differs not from the Spanda principle. The diversity of meditation is due solely to the absolute freedom of Spanda. 62 Sankara is not only the supreme object of devotion; as the static polarity of the absolute, he is the inner reality which holds together its Siva and Sakti outer manifestations. Phenomena are patterns of cognitions projected onto the surface of self-luminous siva-consciousness. There they become apparent, directly revealed to consciousness according to their manifest form. Siva is accordingly symbolized as the ground or surface of awareness, smooth and even like a screen (samabhittitalopama). 64 inscribed on this screen (kutfya) are the countless manifest forms which appear within it rendering it as diverse and beautiful as a fossil ammonite (falagrama). 65 Siva is the sacred ground upon which the cosmic mandalas are drawn, the absolute surface of inscription which bears the mark (cihna) of the universe. Abhinava writes: the variety of this world can only be manifest if the highest lord, who is essentially the pure light of consciousness, exists; just as a surface is necessary for a picture. if external objects were perceived in isolation then, because 'blue' and 'yellow', etc., are self-confined and the perceptions [we have of them] refer to their objects alone and so are insentient, mute and dumb in relation to one another . . . how would it be possible to be aware that an object is variegated? but just as depths and elevations can be represented by lines on a smooth wall, and we perceive [a female figure and think], similarly it is possible to be aware of differences in the variegated (contents of experience) only if all the diverse perceptions are connected together on the one wall of the universal light of consciousness. 66 Siva is the perfect artist who, without need of canvas or brush, paints the world pictures. The instant he imagines it, it appears spontaneously, perfect in every respect. The colours he uses are the varying shades

and gradations of his own Spanda energy and the medium his own consciousness. Intractable disaffect, Mutinous obdurate, Rebellious recalcitrant, Seditious refractory, Insubordinate insurgent, antagonistically antipodean, discrepant incompatible, Contumacious coadjutant, Conflicting converse, Unyielding obdurate, the antagonistic and anachronistic dispositional ties of Anti Brahman are in fact the sine qua non of Brahman only (See Vishwaroopa of Mahavishnu, which contains the polarised positives on RHS and negatives on LHS). The universe is coloured with the dye of its own nature (svabhava) by the power of Siva's consciousness (a//). Rajanaka rama says: homage to him who paints the picture of the three worlds, thereby displaying in full evidence his amazing genius {pratibha)\ to Sambhu who is beautiful with the hundreds of appearances laid out by the brush of his own unique, subtle and pure energy. 68 analogously, at the microcosmic level, all the cognitions and emotions, etc., which make up the individual personality form the outward flow of essentially introverted consciousness. They are specific pulsations (vise$aspanda) or aspects of the universal pulsation (samanyaspanda) of pure t consciousness. at the lower level, within the domain of maya, they represent the play in the fettered soul of the three primary qualities (gunas) or feeling-tones' which permeate to varying degrees his daily experience. These are: 1) sattva — the quality of goodness and luminosity which accompanies blissful experience both aesthetic and spiritual. 2) Rajas — the passion or agitation which oscillates between the extremes of 'light' and 'darkness' and characterises inherently painful experiences. 3) Tamas — the torpor and delusion which accompany states of inertia and ignorance. 69 the liberated soul recognises that these three are the natural and uncreated powers of pure consciousness. for him they are manifest respectively as: 1) Sankara 's power of knowledge (jnana) — the light of consciousness (prakasa); 2) the power of action (kriya) — the reflective awareness of consciousness (vimarsa)\ 3) the power of maya — which does not mean here the world of diversity, but the initial subtle distinction which appears between subject and object in pure consciousness. 70(doctrine of vibration: mark S.G. Dyczkowski) .to "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new

concepts and meanings that could not be included in the old system. at the heart of western metaphysics, for example, derrida finds the opposition between "speech" and "writing." this binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior in derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. The postmodern turn in philosophy: theoretical provocations and normative deficits by Steven best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html. You cannot see beyond consciousness for that is the limit of your storage of information or the basis for understanding further information. Turiya is empty bucket; it is from which prajna the measure of  arises; turiya measures the number of quantum information stores in other three states of waking, dreams, and dreamless deep sleep (Manduka Upanishad; my interpretation for models next paper).let there be path where angels tread: it is a force never born never dies the brahman an anti brahman verily there are parallel universes in (eb) in the ambit of space and time; the society of sinners and the sinned with the grammar of evil; parallel universes contain(e) multiverses in them; all physics, philosophy,  break down(e) in these universes; astronomical gratification ,deprivation, glorification, mortification, projections, axes, dimensions, rotations, foldings, concurrences, motions, shocks, are the axiomatic predications of multiverses and parallel universes; death is when self (the witness consciousness; register of individual general ledger) leaves the body; ostensibly parallel seem to have (e) an objective; parallel universes donot (e) have objectives; their paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common

attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, affection- appreciation, awareness- discernment, emotion-feeling, gut reaction, heart insight, intuition, judgment, keenness, perceptiveness, rationale, sensation, sense, sensitiveness, sensitivity, sentiment, susceptibility, taste, vibes, admiration, aesthetic sense, affection, appraisal, assessment, attraction, awareness, cognizance, commendation, comprehension, enjoyment, esteem, estimation, grasp, high regard, knowledge, liking, love, perceptional realization, recognitional regard, relish, respect, responsiveness, sensibility, sensitiveness, sensitivity, sympathy, understanding, valuation, acknowledgment, airing, baring, , confessional defenselessness, denudation- denunciation, disclosure, display, divulgence-exhibition, hazard, introduction, jeopardy, laying open, liability, manifestation, nakedness, openness, peril, presentation, publicity, revelation, risk, showing, susceptibility , susceptiveness,susceptivity, unfolding, unmasking, unveiling, vulnerability, vulnerableness, action, affection, appreciation and ardor, behavior, capacity, compassion, concern, cultivation and culture, delicacy, discernment and discrimination, emotion, empathy, faculty, fervor, fondness, heat, imagination, impression and intelligence, intensity and intuition, judgment, keenness- palpability, passion, pathos, pity, reaction, refinement, sensibility, sensitivity, sentiment, sentimentality , sharpness, spirit, sympathy, tangibility, taste, tenderness, understanding, warmth, fond memories, hearts and flowers, homesickness, longing, pining, reminiscence, ,schmaltz, sentimentality wistfulness, yearning satiation is only ostensible; their job is to get the job done; their business is to mind other's business ;all laws of physics and finance breakdown; it is antilaw,piscatorial Piratish and bubonic bucaneerishness; there is free transactions between in multiverses; when reflected on the screen of consciousness they appear as extant existential universe. Transactions are astronomical and gratification

deprivation, glorification and mortification incalculable; they have no barriers of region, religion, language which they create for others. theirs objectives are: corruption criminality, debasement debauchery, degradation evil immoral lewdness licentiousness, perversion sinfulness, vice viciousness, vileness wickedness, improbity corruption, atrocity decadence, degeneration degradation, depravity evil, immorality impurity, infamy iniquity, looseness lubricity, perversion profligacy, sinfulness turpitude,vice viciousness, vulgarity wickedness,disgrace bad reputation, abasement abuse, baseness black eye, blemish blur, brand comedown,contempt contumely, corruption culpability, debasement debasing, defamation degradation, derision disbarment, discredit disesteem, disfavor dishonor, disrepute disrespect, humbling humiliationignominy ill repute, infamy ingloriousness, meanness obloquy, odium opprobrium, pollution ,put-down reproach, scandal scorn, slander slight, slur spot, stain stigma, taint tarnish, turpitude venality, enormity Horribleness , abomination atrociousness, atrocity crime and the positive aspectionalities thereof .let there not be any incursions where evil mercenaries in the garb of mendicants thrive with adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatious gainsaying or conformingly pugnacious Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, execrable and goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness . statement for

Mephistophelean mercenaries philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality. Negative energy stabilising wormholes are considered candidates for (er) darkmatter. An analysis of negative energy might be a useful ingredient of the thoughts of travel at speeds more than the velocity of light. Peirce explains that an "imperfect continuum" possesses "topical singularities, or places of lower dimensionality" which break continuity, while "the parts of a perfect continuum have the same dimensionality

as the whole." (CP 4.642) Circular continua without singularities are time and space. (CP 6.210-212) By using math as his lens, Peirce envisioned singularities or lower dimensional holes in the imperfect cosmologic continuum, while maintaining that a probabilistic law of chance exists that entropy decreases **(Wikipedia).** On the side of philosophy**,** quite explicitly, Deleuze provides a critique of Husserl's "sleight-of-hand" change from a static categorical phenomenology to one supposedly more dynamically constitutive of reality. This a clear cut **chagrined circumvention**, **contraventional curbing**, **defeated disgruntlement**, **dissatisfactional downer**, **dragged failure**, **fizzled foil**, **grievance hindrance**, **impedimental irritation**, letdown, nonfulfillment, nonsuccess, obstruction, old one-two, **resent mental setback**, unfulfillment of the earlier theories makes Sartre contend that human existence is a conundrum whereby each of us exists, for as long as we live, within an overall condition of nothingness (nothing-ness)—that ultimately allows for free consciousness. But simultaneously, within our being (in the physical world), we are constrained to make continuous, conscious choices. It is this dichotomy that causes anguish, because choice (subjectivity) represents a limit on freedom within an otherwise unbridled range of thoughts. Subsequently, humans seek to flee our anguish through action-oriented constructs such as escapes, visualizations, or visions (such as dreams) designed to lead us toward some meaningful end, such as necessity, destiny, determinism (God), etc. Thus, in living our lives, we often become unconscious actors—Bourgeois, Feminist, Worker, Party Member, Frenchman, Canadian or American—each doing as we must to fulfill our chosen characters' destinies. However, Sartre contends our conscious choices (leading to often unconscious actions) run counter to our intellectual freedom. Yet we are bound to the conditioned and physical world—in which some form of action is always required. This leads to failed dreams of completion, as Sartre described them, because inevitably we are unable to bridge the void between the purity and spontaneity of thought and all-too constraining action; between the being and the nothingness that inherently coincide in our self. Deleuze insists on the absolute non-resemblance between what conditions (metastable, pre-individual plane of singularities) and what is conditioned

(actualized or the individuated 'thing'). Only then do "the conditions of the true genesis become apparent," Deleuze writes (Logic of Sense 105). Without addressing this issue, Deleuze's borrowings from Simondon, Lautman and structuralism remain untethered. Because of Bowden's decision to completely sidestep Deleuze's overt negotiations with phenomenology, his otherwise sophisticated commentary suggests a pedagogical narrowness. If, the omission of phenomenology strengthens Bowden's focus on the ontological priority of the event over substantial ontologies via the genealogy of figures and movements Deleuze outlines, it also incompletely articulates the problem of the sense-event as Deleuze formulates it. We have lost any idea of whom Deleuze takes to be his primary adversary, against whom he stakes out his position. This worry is most clearly manifested in Bowden's description of Deleuze's ontological prioritizing of events as offering a "transcendental ontology." For Bowden Deleuze's transcendental ontology defines events that "are ontologically prior to substances, 'all the way down'" (82). But Bowden provides a second definition: it is that the priority of events over worldly individuals means that the world of events "is not something external to the conditions of knowledge" (69). **SEAN BOWDEN the Priority of Events: Deleuze's Logic of Sense Sean Bowden, the Priority of Events: Deleuze's Logic of Sense, Edinburgh University Press, 2011, 296pp., $40.00 (hbk), ISBN 9780748643646. Reviewed by David Scott, Coppin State University Weltschmerz, agony, apprehension, blues, depression, dread, mid-life crisis, misgiving, nervousness, uneasiness, angst were expressed with such effulgent words…… aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement, argument, articulational assertion, asseverational choice of words, commentary and communication, declaration and definition, delivery, diction, elucidation, emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation about the human predicament is brought out in Being and Nothingness. Being and Nothingness offers a critique of Sigmund Freud's theory of the unconscious, based on the claim that consciousness is essentially self-conscious. Sartre also argues that**

**Freud's theory of repression is internally flawed. According to Sartre, in his clinical work, Freud encountered patients who seemed to embody a particular kind of paradox—they appeared to both know and not know the same thing. In response, Freud postulated the existence of the unconscious, which contains the "truth" of the traumas underlying the patients' behavior. This "truth" is actively repressed, which is made evident by the patients' resistance to its revelation during analysis. Yet what does the resisting if the patients are unaware of what they are repressing? Sartre finds the answer in what Freud calls the "censor". "The only level on which we can locate the refusal of the subject," Sartre writes, "is that of the censor.** The Holographic Principle, as conceived in current physics, applies to fields, and perhaps, even to more elementary entities, called strings and branes. These more elementary entities remain hypothetical at the time of this writing. There are many layers between the level of fields and that of neurons. We know that neurons are surrounded by a surface phospholipid membrane, which supports an electrochemical process that is fundamentally necessary to human experience. As the information supported by the membrane surface according to the Holographic Principle is proportional to its area, we should expect to find that the neural correlates of consciousness most clearly associated with those parts of the neuron that have the have a high proportion of neuronal surface area, and, in particular, receptive surface area. Dendrites account for an average of about 90% of the neuron is receptive surface area (Wong, 2002), and so would be expected to be the most important neural structures in information processes. We find this to be true **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science (Wikipedia).** Discordant-discrepant, dissonant-hard-line, implacable, incompatible, incongruous- inconsistent, inexorable-inflexible, inharmonious-intransigent, opposed, reluctant, unappeasable, uncompromising, unfriendly some of these expatiations and enumerations might sound, a more generalised interpretations is given by a Vedantin. Look at this arresting passage: The teacher continued to be silent. When addressed a second and third time he said: "I am teaching, but you do not follow. The Self is silence." The undetermined and unthinkable character of the Brahman is

a consequence of the absolute's eternal and immutable nature. To concede the existence of a real universe is, from the Vedantin's point of view, to posit the existence of a reality apart from the Brahman. Nor can we simply identify a real universe with the absolute unless we are prepared to compromise its unchanging, absolute status. The criterion of authenticity is immutability. Reality never changes; only that which is less than real can appear to do so. Reality is constant in the midst of change. What this means essentially is that there is change although nothing changes. This impossible situation is reflected in the ultimate impossibility of change itself. That which does not exist prior to its changing and at the end, after it has changed, must be equally non-existent between these two moments. Although the world of change appears to be real, it cannot be so. Change, according to the Vedantin, presupposes a loss of identity. Reality cannot suffer transformation; if it were to do so, it would become something else and the real would be deprived of its reality. The immortal can never become mortal, nor can the mortal become immortal. The ultimate nature of anything cannot change. Change of any sort is merely apparent (vivarta); the world of change and becoming is a false super- imposition (adhyaropa, adhyasa) on the absolute. In cosmic terms, the mistake (bhranti) consists of the supposition that the real Brahman is the unreal universe and the unreal universe is the real Brahman. In microcosmic terms, it is the mistake of falsely conceiving the body, mind or even one's personality to be the Self. In the same way as the image of a snake is falsely superimposed on a rope, similarly the universe is falsely projected onto the real substratum, the Brahman. Ignorance is not merely a personal lack of knowledge, but a cosmic principle. As such it is called "Maya," the indefinable factor (anirvacaniya) that brings this mistake in identity about. The reality status of this cosmic illusion is also indefinable: on the one hand it is not Brahman, the sole reality; on the other hand it is not absolutely non-existent like a hare's horn or the son of a barren woman. Brahman is the source of world appearances only in the sense of being their unconditioned ground or essential nature. The universe is false not because it has no nature of its own but because it does have one. Just as the illusion of a snake disappears when one sees that it is nothing but a rope, similarly cancellation (badha) of the empirically real occurs

when the absolute reality of the Brahman is realised. Thus, according to Vedanta, appearance implies the real, while the real need not imply appearance. To appear is essentially to appear in place of the real, but to be real is not necessarily to appear. All things -exist because the absolute exists. It is their Being. Thus the very existence of phenomena implies their non-existence as independent realities.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^\infty \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9 \left( U_9 - V_9 \int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9 \right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9 \int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$then given the constraints of $f(T_9)$, $(A1.5)$*can never be satisfied Co*nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 -$$

V 9M 9–]**e**–*T*9λ 9=0

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9 Z_9 \sin(Z_9 T_9) - (A_9 M_9 + Y_9) + (A_9 M_9 + X_9)\cos(Z_9 T_9) = 0$$

$$(A_9 M_9) Z_9 - M_9 Z_9 \cos(Z_9 T_9) + (A_9 M_9 - X_9)\sin(Z_9 T_9) = 0$$

Quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High

Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) **Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP** This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) **Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP** In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference: Journal of Physics: Conference Series 70 **Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows:** System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol

Where $A_9 = U_9 - V_9$Using (A2.3), substitute for $\boldsymbol{Z}_9^2$ and $Z_9$in (A2.2):

$W_9^2 (X_9^2 - 3 M_9 A_9 X_9 - M_9^2 (X_9 - Y_9) + 2 M_9^3 A_9 + 2 A_9^2 M_9^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_9$=1-cos ($Z_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (X_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ X_9 - Y_9/2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \leq W_9 \leq 2$ *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (X_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\,M_9)^2\,(2A_9 M_9 - (X_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9 M_9^2 - 3M_9 X_9\right) + W_9\left(3M_9 A_9^2 - 2A_9 X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9+Y_9\right)$$

$$=2M_9^2W_9^2(M_9+\mathrm{A}_9)+M_9A_9^2W_9+\left(2A_9M_9-X_9+Y_9\right)(A_9W_9+M_9W_9^2$$

$$+(2-W_9)((2M_9X_9+M_9Y_9)W_9+A_9(X_9+Y_9)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range $0 \le W_9 \le 2$.Defining

$Q_9\left(W_9,A_9\right)=a_9W_9^2+c_9$ *where* $a_9\,b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

To Heidegger writings we shall compose the model. Following is not still written model able form. For one of the models please see any one of the papers under the appellation: Ïn Excited Reveries

In his later writings Heidegger drops the term 'phenomenon' in preference for the verbal form 'phainesthai' in order to emphasize even more the actuality or presentational property of Being. Explaining this new form of the term he writes: "Being disclosed itself to the ancient Greeks as 'physis'. The etymological roots 'phy-' and 'pha-' designate the same thing: 'phyein', the rising-up or upsurge which resides within itself as 'phainesthai', lighting-up, self-showing, coming-out, appearing-forth." Heidegger contrasted his notion of phenomenon with semblance (Schein) and with appearing (Erscheinung). In the case of semblance a thing can show itself as that which it is not, as when fool's gold shows itself to be gold. The ancients always allied semblance with non-being. Heidegger points out, however, that semblances are grounded in

showings, and so does Abhinava. Both Heidegger and Abhinava consequently maintain that all semblances have a real basis and are to be treated as instances of phenomena along with the so-called real showing or manifestation of non-deceptive objects. So Heidegger states that: 'how- ever much seeming, just that much being'. Thus self-showing or appearing defines Being as phenomenon, but this definition of Being is as yet incomplete. Being is not only self-showing but 'logos' which Heidegger explains means 'discourse' (Rede) in the sense of 'apophansis': 'letting-be-seen'. Phenomenology, which according to Heidegger is the only correct study of Being, means 'letting-be-seen-that-which- shows-itself. This is true of Saiva Paramadvaita as well. The reality of the world demands recognition; we are forced to accept the direct presentation of the fact of our daily experience. As Abhinava says: "if practical life, which is useful to all persons at all times, places and conditions were not real, then there would be nothing left which could be said to be real" A thousand proofs could not make 'blue' other than the colour blue. The reality of whatever appears in consciousness cannot be denied. Objects appear; they do not cease to do so by a mere emphatic denial. The manifestation of an entity in its own specific form is a fact at one level of consciousness; it is real. The appearing of the same entity in the same form but recognised to be a direct representation of the absolute is also a fact, but at another level of consciousness. It is no more or less real than the first. 'As is the state of consciousness, so is the experience,' says Abhinava. Although the nature of the absolute is discovered at a higher level of consciousness, The Doctrine of Vibration nonetheless it presents itself to us directly in the specific form in which we perceive things; otherwise there would be no way in which we could penetrate from the level of appearing to that of its source and basis. Abhinava writes: Real is the entity iyastu) that appears in the moment of direct perception (sak$atkara), that is to say, within our experience of it. Once its own specific form has been clearly determined one should, with effort, induce it to penetrate into its pure conscious nature. All things are known to be just as they present themselves. The concrete actuality of

being known (pramiti), irrespective of content, is itself the vibrant (spanda) actuality of the absolute. Liberating knowledge is gained not by going beyond appearances but by attending closely to them. "The secret," Mahesvarananda says, "is that liberation while alive (jivanmukti) is the profound contemplation of Maya's nature." No ontological distinction can be drawn between the absolute and its manifestations because both are an appearing (dbhasa), the latter of diversity and the former of 'the true light of consciousness which is beyond Maya and is the category Siva'. Those who have attained the category of Pure Knowledge above Maya and have thus gone beyond the category of Maya, see the entire universe as the light of consciousness . . . Just as the markings [on a feather] are nothing apart from the feather, the feather [is nothing apart from] them, similarly, when the light of consciousness is manifest, the whole group of phenomena is manifest as the light of consciousness itself. Within the sphere of Maya, every entity's 'own nature' (svabhdva) corresponds to its specific manifest form. Accordingly it is defined as that which distinguishes it from all else and from which it never deviates. Above the sphere of Maya, that is, above the level of objectivity, is the domain of the subject. At this level, everything is realised to be part of the fullness of the experience r and hence no longer bound by the conditions which impinge on the object. Here the part is discovered to be the whole, that is, consciousness in toto. In this sphere beyond relative distinctions, the yogi realises that (all) the categories of existence are present in every single category. m The yogi experiences every individual particular as the sum total of everything else. He recognises that all things have one nature and that every particular is all things. This is the 'essence' (sard) or co-extensive unity (samarasya) of all things. Integral Monism of Kashmiri SaivismWe have established that reality is manifest according to how [and the degree in which] the freedom of consciousness reveals it and that [this freedom] is the womb of all forms. Just as 'sweetness' is present in its entirety in every atom of the sugarcane, so each and every atom [of the universe] bears within itself the emanation of all things. This is

the level of consciousness in which the absolute reflects on itself realising to its eternal delight and astonishment (camatkara) its own integral nature. The reality of the world of diversity is not denied, but experienced in a new mode of awareness free of time and space in the eternal omnipresence of the Here and Now. [Phenomenal forms of awareness] such as 'this [exists]', born of the colouring [imparted to the absolute] by the limitations engendered by the diversifying power of time {kalakalana) also emanate within the Supreme Principle. There [at that level], Fullness {purriata) is the one nature [of all things] and so everything is omnipresent; otherwise, associated with division (khancjlana), the Fullness [of the absolute] would not be full. The content of absolute consciousness consists of diverse appearings (abhdsa) which, because they are manifest through it in this way, do not compromise the wholeness of consciousness. Everything we perceive is a momentary collocation of a number of such manifestations which combine together like 4 a row of altar lamps' (dipavalT) to form the single radiant picture of the universe. The individual objects which constitute the universe are specific collocations of such 'atomic' appearings. Together they form a single unified particular which appears according to its own defining features (svalaksand). A jar, for example, consists of a number of appearances such as 'round', 4 fat\ 'earthen', 'red', etc., which together discharge a single function (arthakriya), in this case, that of carrying the appearance 'water'. They unite with each other much as the scattered rays of a lamp come together when focused, or as the various currents of the sea together give rise to waves. Atomic appearings can combine in any number of ways, provided that they are not contrary to one another as established by the dictates of natural law (niyati). An appearance of 'form', for example, cannot combine with that of 'air'. Insofar as they share a common basis (samanyadhikaranya), a given cluster of appearances appears as a single whole. This common basis is the most prominent member of the group; the appearance 'jar' is such in the example quoted above. Any one appearance in a cluster may assume a more important or subordinate role.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_9 - b_9 < 0$$

And

$$b_9^2 > 4a_9c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9 in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\ X_9$ ,

$Let\ Y_9 = y_9\ X_9$And try the substitution $A_9 = X_9/(2 + A_9\ y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + n_9 y_9)$ we have

$$W_9^2 y_9\left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^2\left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9\,X_9\left[\frac{X_9\left(1+n_9y_9\right)}{M_9^2\left(2+n_9y_9\right)^2}+2\,y_9\left(n_9-1\right)+\frac{X_9^2\left(9+y_9\right)}{M_9^2\left(2+n_9y_9\right)}\right]=0.\ (A2.12)$$

Since $n_9$ and $A_9$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9=\left(2+n_9y_9\right),$ *the condition*$(A2.99?)$*can be written as*

$$\left[\frac{X_9\left(E_9-9\right)}{E_9^2}\right]^2\left[\frac{X_9^2}{M_9^2}-4\,y_9\,E_9^2\frac{X_9}{M_9^2}+4y_9\left(n_9-2\right)E_9^3\right]\ \square 0.$$

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9/W_9^2)$ can be expressed as

$$2\,E_9\left[y_9E_9+\left(y_9\,E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}\right]>\frac{X_9}{M_9^2}>2E_9\left[y_9E_9-y_9E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}]$$

dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique

for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) **Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP** Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) **Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP** Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a

cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) **Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP** Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service

placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) **Dynamic Service Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP** Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by $H*(X,(\wedge*TX)\otimes g)$, where X is the target hyperk\"ahler manifold. We further give a mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) **A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van**

**Raamsdonk1 and Nathan Seiberg2** study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. **Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP** Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy **The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP** Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the

present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. **ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP** Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition. Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb) the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537

© 2016 American Physical Society **Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP** Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant ($\Delta\alpha/\alpha$) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and $\Delta\alpha/\alpha$ to (e&eb) this conclusion. However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts $z\lesssim 3$, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the $\Lambda$ cold dark matter model value ($j\Lambda CDM=1$); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in $\alpha(z)$, linked to (e&eb) a linear coupling ($\zeta$) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order $|\zeta|\sim 10-4$, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in $\Delta\alpha/\alpha$ (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society **The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP** Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation

rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void

Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems, authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society **Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP** Authors measure the star formation properties of (e) two large samples of

galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the **star formation** rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a **function of environment**. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]**.** Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size $\sqrt{n} \times \sqrt{n}$? The problem is that $\sqrt{n}$ time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time $O(\sqrt{n})$ for $d \geq 3$, or $O(\sqrt{n}\log 5/2n)$ for $d = 2$. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search

tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ∗This work was mostly done while the author was a **PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship and by CIAR**. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 **Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP** Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can

lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 **Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP** Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. **Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP** Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e)

sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) **Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP** The relationship between the foundations of mathematics and (e&eb) physics is a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb) flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling

field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) **Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP** The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar

and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects:

Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph]

(or arXiv: 1503.05600v1 [math-ph] for this version) **Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP** This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs

mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) **Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP** Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and

derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to **the fine structure constant**. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) **Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP** Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor goes to 0 as the (e&eb) time goes to 0, the

big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph] (or arXiv: 1306.4613v1 [quant-ph] for this version) **Effects of gauge theory based number scaling on geometry Paul Benioff STOP** In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x,

outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph]

(or arXiv: 1211.3381v2 [quant-ph] for this version) **Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP** Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value,

\int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances, the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) **Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP** The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include

(e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) **Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP** Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) **New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP** This work is based on (e) the field of reference frames based on (e) quantum

representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects:Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc) Journal reference:Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) **A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul B**

**ABSOLUTE STABILITY**

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 ? $W_9$ ? 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9-9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2-2y_9\frac{(E_9+n_9^2y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right)-4y_9(E_9-n_9)$$ (A2.14)

And has the solution

$$\frac{X_9}{M_9^2}>\frac{E_9}{(E_9-9)}[y_9(E_9+n_9^2y_9)+[y_9^2(E_9+n_9^2y_9)^2+4y_9(E_9-9)(E_9-n_9)]^{9/2}]$$

Given $E_9>n_9$this lower limit to $X_9/M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9/M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9<n_9$ $(ie.,2A_9M_9^2<X_9-Y_9)$then (A2.9)and (A2.8)show that $Q_9$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2}>2E_9[y_9E_9-(y_9E_9(9+y_9)(E_9-n_9)]^{9/2}]$$

And

$$E_9>n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2/2M_9(X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$ and $X_9$ and $Y_9$ by $W_9^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

**Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability

of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e) long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007 STOP. David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer P** examine the possibility of coherent reversible information transfer between solid-state superconducting qubits and (e&eb) ensembles of ultracold atoms. Strong coupling between these systems is mediated by (e&eb) a microwave transmission line resonator that interacts near resonantly with (e&eb) the atoms via (e&eb) their optically excited Rydberg states. Solid-state qubits can then be used to (e) implement rapid quantum logic gates, while (e) collective metastable states of (e) the atoms can be employed for (e) long-term storage and (e&eb) optical readout of quantum information. **Reversible state transfer between superconducting qubits and (e&eb) atomic ensembles David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer Phys. Rev. A 79, 040304(R) – Published 14 April 2009 STOP.** Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures have been proposed. An

attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. **Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier** demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in (eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. **Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP.** In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown

to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , **Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally** frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. **Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg** propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states

into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. **Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011 http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan** proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 **Qubit transfer between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller** propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. **Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys.**

**Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich** report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. **Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi** investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. **Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev** investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer.

Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. **L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP.** In recent years **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P have** investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P** apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys. Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. **Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S.**

**Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP.** Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert **Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe** realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.**220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2** [quant-ph] for this version) **Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP.** Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. **Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker** describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue

from over (e&eb) a thousand plant and animal samples. **NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP.** Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this behavior, enhancing and accelerating (eb, eb+) the growth of correlations. **Michael Kastner** gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. **Michael Kastner** tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) **Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP.** A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. **Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E. Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91,**

**075435 – Published 26 February 2015 STOP.** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still

problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to

(e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP** Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that

(eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035

**Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime**

**C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013**

**STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and

cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP** Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex

Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) **A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen**

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9 and$ separate (A3.1)into real and imaginary parts

$$Z_9^2 + \left(M_9 + V_9\right) Z_9 \sin\left(Z_9 T_9\right) - \left(U_9 M_9 + Y_9\right) + (M_9\left(U_9 + V_9\right) - X_9)$$

$$\mathrm{X}\cos\left(Z_9 T_9\right) - 2M_9 V_9 cos^2\left(Z_9 T_9\right) + M_9 V_9 = 0$$

$$(M_9 + U_9)\, Z_9 - (M_9 + V_9)\, Z_9\ cos\left(Z_9 T_9\right) + \left(M_9\left(U_9 + V_9\right) - X\right)$$

$$\mathrm{X}\sin\left(\left(Z_9 T_9\right) - M_9 V_9 \sin\left(Z_9 T_9\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 =$ $1 - cos\, Z_9\ T_9$and $A_9 = U_9 - V_9$we obtain

$$4\, M_9^2 V_9^2 W_9^4 + 2\, M_9 V_9\ W_9^3[(M_9 - V_9)^2 + A_9\left(M_9 + V_9\right) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2\left[M_9 - V_9\right)^2\left(2\, A_9\ M_9 + Y_9 - X_9\right) + \left(2A_9\ M_9 - X_9\right)^2 - A_9\left(M_9 - V_9\right)\left(2A_9\ M_9 - X_9\right)$$

$$+ M_9 V_9\,(A_9^2 + 4\left(X_9 + Y_9\right))]$$

$$+ W_9\left[A_9^2 + 2X_9\right)\left(2\, A_9\ M_9 - X_9\right) + 2A_9 Y_9\,(M_9 + V_9) - M_9 A_9^3] + A_9^2\left(X_9\ M_9 - Y_9\right) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9\left(2, U_9 V_9\right) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting cos $(Z_9 T_9) = -1\ in\ (A3.3)$

We have applied the model to Kashmiri Shaivism. Whose elaborate rituals is enough to make a indivual raise from individual consciousness to cosmic consciousness which we have called Cosmic Consciousness.

It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) **Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP** The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic

Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** The group of unimodular transformations on the roots of polynomials over GF(2) is considered, and those polynomials with (e&eb) symmetries in (eb) the unimodular group are identified. The cross-correlation function between (e&eb) two maximum-length linear shift register sequences of (e) the same degree is shown to be (=) computable as an explicit linear transformation, in (eb) matrix form, on (e&eb) either one of the sequences, regarded as (=) a vector. The underlying vector space is (=) the "cyclotomic algebra," generated by (e) the cyclotomic cosets, or by what Gauss termed the **"periods" of the cyclotomic equation**. A variety of numerical examples are worked in detail. **Information Sciences Volume 1, Issue 1, December 1968, Pages 87-109 Theory of transformation groups of polynomials over GF(2) with applications to linear shift register sequences Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more link. Solomon W. Golomb** stop Authors study the P−V criticality and phase transition in (eb) the extended phase space of charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of (e) the system and its conjugate quantity is (=) the thermodynamic volume of the black hole. The black holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb) no P−V criticality and phase transition appear, while (e) for the black holes with a spherical horizon, even when the charge of the black hole is absent, t(eb) he P−V criticality and the small black hole/large black hole phase transition will appear, but it happens only in d=5 dimensions; when the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to have (e) the P−V criticality

and (e&eb)the small black hole/large black hole phase transition, there exists (ebb) an upper bound for the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the **Gauss-Bonnet coefficient** and Q is (=) the charge of the black hole. Authors calculate the critical exponents at the critical point and find that (eb) for all cases, (eb) they are the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09(2013)005 Cite as: arXiv:1306.6233 [gr-qc] (or arXiv:1306.6233v4 [gr-qc] for this version) **P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in IR but with (e&eb) improved UV behaviors. Here, authors study an IR modification which breaks ``softly'' (e) the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Authors obtain the black hole and cosmological solutions for **``arbitrary'' cosmological constant** that represent (eb) the analogs of the standard Schwarzschild-(A)dS solutions which can be asymptotically (A)dS as well as flat and authors discuss their thermodynamical properties. They also obtain solutions for (e) FRW metric with (e&eb) an arbitrary cosmological constant. Authors study its implication to (e&eb) the dark energy and find that it seems to be consistent with (e&eb) current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP.** The Einstein equations with (e&eb) a negative cosmological constant admit (e) black hole solutions which are asymptotic to (e) anti-de Sitter space. Like black holes in asymptotically flat space, these solutions have (e) thermodynamic properties including (e) a characteristic temperature and intrinsic entropy equal to (=) one quarter of the area of the event horizon in Planck units. There are however some important differences from (e) the asymptotically flat case. A black hole in anti-de Sitter space has (e) a minimum temperature which occurs when (e) its size is of the order of the characteristic radius of (e) the anti-de Sitter space. For larger black holes the red-shifted

temperature measured at (eb) infinity is greater. This means that such black holes have (e) positive specific heat and can be in stable equilibrium with (e&eb) thermal radiation at a fixed temperature. It also implies (eb) that the canonical ensemble exists for asymptotically anti-de Sitter space, unlike the case for asymptotically flat space. One can also consider the microcanonical ensemble. One can avoid (e) the problem that arises in asymptotically flat space of (e) having to put the system in (eb) a box with unphysical perfectly reflecting walls because (e) the gravitational potential of anti-de Sitter space acts as a box of finite volume. Communicated by A. Jaffe **Communications in Mathematical Physics December 1983, Volume 87, Issue 4, pp 577-588 Thermodynamics of black holes in anti-de Sitter space S. W. Hawking, Don N. Page STOP** In previous work authors introduced the concept of black-hole entropy, which they identified with (e&eb) the surface area of the black hole in question expressed in units of (e) the Planck length squared. Authors suggested that the appropriate generalization of the second law for a region containing a black hole is (=) that the black-hole entropy plus (+) the common entropy in the black-hole exterior never decreases. Here authors establish the validity of this law for the infall of an entropy-bearing system into (e&eb) a much larger and more massive generic stationary black hole. To do this authors determine (eb) a general lower bound for the increase in black-hole entropy, and an upper bound for the entropy of the system, while allowing for quantum effects at each stage. In passing authors show that the generalized second law is a statistical law which becomes (=) over-whelmingly probable in the limit of a macroscopic system. Authors also consider briefly more general situations. Finally, we give two simple examples of predictions made by the generalized second law for (e) black-hole formation processes. Received 17 September 1973 DOI:http://dx.doi.org/10.1103/PhysRevD.9.3292 **Generalized second law of thermodynamics in black-hole physics Jacob D. Bekenstein Phys. Rev. D 9, 3292 – Published 15 June 1974 STOP** The Lovelock gravity consists of (e) the dimensionally extended Euler densities. The geometry and horizon structure of (e) black hole solutions could be quite complicated in (eb) this gravity; however, authors find that some thermodynamic quantities

of the black holes like the mass, Hawking temperature and entropy have (e) simple forms expressed in terms of (e) horizon radius. The case with black hole horizon being a Ricci flat hypersurface is (=) particularly simple. In that case the black holes are (=) always thermodynamically stable with (e&eb) a positive heat capacity and their entropy still obeys (e&eb) the area formula, which is no longer valid for (e) black holes with positive or negative constant curvature horizon hypersurface. In addition, for black holes in (eb) the gravity theory of Ricci scalar plus a 2n-dimensional Euler density with (e&eb) a positive coefficient, thermodynamically stable small black holes always exist in (eb) D=2n+1 dimensions, which are absent in (eb) the case without (e) the Euler density term, while (e) the thermodynamic properties of the black hole solutions with (e&eb) the Euler density term are qualitatively similar to those of black holes without (e) the Euler density term as D>2n+1. Copyright © 2004 Published by Elsevier B.V. **Physics Letters B Volume 582, Issues 3–4, 4 March 2004, Pages 237–242 A note on thermodynamics of black holes in Lovelock gravity Rong-Gen CAI STOP** There are a number of similarities between black-hole physics and (e, eb, =) thermodynamics. Most striking is the similarity in the behaviors of black-hole area and (e&eb) of entropy: Both quantities tend to increase (eb+) irreversibly. In this paper we make this similarity the basis of (e) a thermodynamic approach to black-hole physics. After a brief review of the elements of the theory of information, authors discuss black-hole physics from the point of view of information theory. Authors show that it is natural to introduce (eb) the concept of black-hole entropy as (=) the measure of information about (e&eb) a black-hole interior which is (=) inaccessible to an exterior observer. Considerations of simplicity and consistency, and dimensional arguments indicate (eb) that the black-hole entropy is equal to (=) the ratio of the black-hole area to the square of the Planck length times a dimensionless constant of (e) order unity. A different approach making use of the specific properties of Kerr black holes and of concepts from information theory leads to (eb) the same conclusion, and suggests (eb) a definite value for the constant. The physical content of (e) the concept of black-hole entropy derives from (e) the following generalized version of the second law: When common entropy goes down a

black hole, (eb) the common entropy in (eb) the black-hole exterior plus the black-hole entropy never decreases. The validity of this version of the second law is supported by (e) an argument from information theory as well as by several examples. Received 2 November 1972 DOI:http://dx.doi.org/10.1103/PhysRevD.7.2333 **Black Holes and Entropy Jacob D. Bekenstein Phys. Rev. D 7, 2333 – Published 15 April 1973 STOP** The effects of power-law Maxwell field on the Van der Waals like phase transition of higher-dimensional dilaton black holes are probed in detail. It is shown that the Smarr relation gains corrections due to (e) the effects of (e) both the power-law Maxwell field and (e&eb) the dilaton field while (e) thermodynamic volume is exactly the same as (=) that of Einstein-Maxwell-dilaton black holes. Authors successfully derive the analytic solutions of (e) critical point and carry out (e&eb) some check to ensure that these critical quantities are (=) positive. It is shown that the constraint on the parameters turns out to be (=) 0<\alpha^2<1, which is tighter than that in (eb) the non-extended phase space. It is also shown (eb) that these critical quantities and (e&eb) the ratio P_cv_c/T_c are affected by (e&eb) the power-law Maxwell field. Moreover, critical exponents are found to coincide with those of other AdS black holes, showing (eb) the powerful influence of mean field theory. Comments: references added Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1601.05500 [gr-qc] (or arXiv:1601.05500v2 [gr-qc] for this version **Effects of power-law Maxwell field on the Van der Waals like phase transition of higher dimensional dilaton black holes Jie-Xiong Mo, Gu-Qiang Li, Xiao-Bao Xu STOP** A consistent and unified picture for (e) critical phenomena of charged AdS black holes in f(R) gravity is drawn in this paper. Firstly, we investigate the phase transition in canonical ensemble. Authors derive the explicit solutions corresponding to (e&eb) the divergence of CQ. The two solutions merge into (e&eb) one when (e) the condition Qc=−13R0−−−√ is satisfied. The curve of specific heat for Q<Qc has two divergent points and can be divided into (e&eb) three regions. Both the large radius region and the small radius region are (=) thermodynamically stable with (e&eb) positive specific heat while (e) the medium radius region is unstable with (e&eb) negative specific heat. However, when Q>Qc, the specific heat is

always positive, implying (eb) the black holes are locally stable and no phase transition will take place. Secondly, both the T−r+ curve and T−S curve f(R) AdS black holes are investigated and they exhibit Van der Vaals like behavior as the P−v curve in the former research. Critical physical quantities are obtained and they are consistent with (=) those derived from the specific heat analysis. Authors carry out numerical check of Maxwell equal area law for the cases Q=0.2Qc, 0.4Qc, 0.6Qc, 0.8Qc. The relative errors are (=) amazingly small and can be negligible. So the Maxwell equal area law holds for (e) T−S curve of f(R) black holes. Thirdly, authors establish geometrothermodynamics for (e) f(R) AdS black hole to examine the phase structure. It is shown that the Legendre invariant scalar curvature R would diverge exactly where (e) the specific heat diverges. To summarize, the above three perspectives are consistent with (e&eb) each other, thus providing (eb) a unified picture which deepens the understanding of (e) critical phenomena of (e) f(R) AdS black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1602.01251 [gr-qc] (or arXiv:1602.01251v1 [gr-qc] for this version) **A consistent and unified picture for critical phenomena of f(R) AdS black holes Jie-Xiong Mo, Gu-Qiang Li, Yu-Cheng Wu STOP** The state equation of the charged AdS black hole is reviewed in (e&eb) the T−r plane. Thinking of the phase transition, the T−S, P−V, P−v graphs are plotted and then the equal area law is used (e) in the three cases to get (eb) the phase transition point (P, T). The analytical phase transition point relations for (e) P-T of charged AdS black hole has been obtained successfully. By comparing the three results, authors find that the equal area law possibly cannot (e) be used directly for P−v plane. According to the T−S, P−V results, they plot (e&eb) the P−T−Q graph and find that for a highly charged black hole a very low temperature condition is required for (e) the phase transition. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI:10.1140/epjc/s10052-015-3641-0 Cite as: arXiv: 1503.07658 [gr-qc] (or arXiv:1503.07658v1 [gr-qc] for this version) **A note on Maxwell's equal area law for black hole phase transitionShan-Quan Lan, Jie-Xiong Mo, Wen-Biao Liu STOP** Recently, extended phase space thermodynamics of Lovelock AdS black holes has been of great interest. To provide

insight from (e) a different perspective and gain (e) a unified phase transition picture, non-extended phase space thermodynamics of (e) (n+1)-dimensional charged topological Lovelock AdS black holes is (=) investigated detailedly in (eb) the grand canonical ensemble. Specifically, the specific heat at (eb) constant electric potential is calculated and phase transition in the grand canonical ensemble is discussed. To probe the impact of the various parameters, they utilize (e) the control variate method and solve (e) the phase transition condition equation numerically for the case k=1, −1. There are two critical points for (e) the case n=6, k=1 while there is only one for other cases. For k=0, there exists no phase transition point. To figure out the nature of phase transition in the grand canonical ensemble, authors carry out an analytic check of the analog form of (e) Ehrenfest equations proposed by Banerjee et al. It is shown that Lovelock AdS black holes in (eb) the grand canonical ensemble undergo (e&eb) a second order phase transition. To examine the phase structure in (eb) the grand canonical ensemble, authors utilize the thermodynamic geometry method and calculate both the Weinhold metric and Ruppeiner metric. It is shown that for both analytic and graphical results that the divergence structure of (e) the Ruppeiner scalar curvature coincides with (=) that of the specific heat. Research provides (eb) one more example that Ruppeiner metric serves as (=) a wonderful tool to (e) probe the phase structures of (e) black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C75 (2015) 5, 211 DOI:10.1140/epjc/s10052-015-3434-5 Cite as: arXiv: 1503.01956 [gr-qc] (or arXiv: 1503.01956v1 [gr-qc] for this version **Non-extended phase space thermodynamics of Lovelock AdS black holes in grand canonical ensemble Jie-Xiong Mo, Wen-Biao Liu STOP** To provide an analytic verification of the nature of phase transition at the critical point of P−V criticality, the original expressions of Ehrenfest equations have been introduced directly. By treating the cosmological constant and (e&eb) its conjugate quantity as (=) thermodynamic pressure and volume respectively, authors carry out analytical check of (e) classical Ehrenfest equations. To show that approach is universal, authors investigate not only higher-dimensional charged AdS black holes (e&eb), but also rotating AdS black holes. Not only are the examples

of Einstein gravity shown, but also the example of modified gravity is presented for (e) Gauss-Bonnet AdS black holes. The specific heat at constant pressure CP, the volume expansion coefficient α and the isothermal compressibility coefficient κT are found to diverge exactly at the critical point. It has been verified that both Ehrenfest equations hold at (eb) the critical point of P−V criticality in (eb) the extended phase spaces of (e) AdS black holes. So the nature of the critical point of P−V criticality of AdS black holes has been demonstrated analytically to be (=) a second-order phase transition. These results are consistent with (e&eb) the nature of liquid-gas phase transition at the critical point. In this sense, extant research would deepen the understanding of (e) the relations of AdS black holes and (e&eb) liquid-gas systems. Moreover, successful approaches to introduce (eb) the original expressions of Erhenfest equations directly into (e&eb) black hole phase transition research demonstrate again (eb) that black hole thermodynamics is closely related (e&eb) to classical thermodynamics, which allows (eb) us to borrow techniques from (e) classical thermodynamics to investigate the thermodynamics of black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 89, 084057 (2014) DOI: 10.1103/PhysRevD.89.084057 Cite as:arXiv: 1404.3872 [gr-qc] (or arXiv:1404.3872v1 [gr-qc] for this version) **Ehrenfest scheme for P−V criticality of higher dimensional charged black holes, rotating black holes and Gauss-Bonnet AdS black holes Jie-Xiong Mo, Wen-Biao Liu STOP** Aiming at a unified phase transition picture of (e) the charged topological black hole in (eb) Ho\v{r}ava-Lifshitz gravity, authors investigate this issue not only in (eb) canonical ensemble with (e&eb) the fixed charge case but also in grand-canonical ensemble with (e&eb) the fixed potential case. Authors firstly perform the standard analysis of the specific heat,(e&eb) the free energy and the Gibbs potential, and (e&eb) then study its geometrothermodynamics. It is shown that the local phase transition points (eb) not only witness the divergence of (e) the specific heat, but also witness the minimum temperature and the maximum free energy or Gibbs potential. They also witness the divergence of (e) the corresponding thermodynamic scalar curvature. No matter which ensemble is chosen, the metric constructed can

successfully produce (eb) the behavior of the thermodynamic interaction and (e&eb) phase transition structure while other metrics failed to (e) predict the phase transition point of (e) the charged topological black hole in former literature. In grand-canonical ensemble, authors have discovered the phase transition which has not been reported before. It is similar to (e) the canonical ensemble in which the phase transition only takes place when k=−1. But it also has its unique characteristics that the location of the phase transition point depends on (eb) the value of potential, which is different from (e) the canonical ensemble where (e) the phase transition point is (=) independent of the parameters. After an analytical check of Ehrenfest scheme, authors find that the new phase transition is (=) a second order one. It is also found that the thermodynamics of (e) the black hole in Horava-Lifshitz gravity is quite different from that in (eb) Einstein gravity. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP1310(2013)056 DOI:10.1007/JHEP10(2013)056 Cite as: arXiv:1404.2497 [gr-qc] (or arXiv:1404.2497v1 [gr-qc] for this version) **A unified phase transition picture of the charged topological black hole in Horava-Lifshitz gravity Jie-Xiong Mo, Xiao-Xiong Zeng, Gu-Qiang Li, Xin Jiang, Wen-Biao Liu STOP** To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case $\beta\to\infty$ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness condition of (e) entropy. It is shown

that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure authors have discovered. It is shown that Legendre invariant thermodynamic scalar

curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws, leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions, geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged AdS black holes by considering the case of the RN-AdS black holes

surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase transition. For the case $\omega_q=-2/3$ authors derive analytic expressions of (e) critical physical quantities, while (e) for cases $\omega_q\neq-2/3$ appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this

version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=) chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy

Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed

quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions, while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma?**

All this consists of the act of consciousness alone", says Vasubandhu, "because unreal entities appear, just as a man with defective vision sees unreal hair or a moon, etc."He points to dreams as examples of purely subjective constructs which appear to be objective realities. The apparent reality dreams possess is not derived from any concrete, objective world, but merely from the idea of objectivity. While the Yogacara does not say that an idea has, for example, spatial attributes, it does have a form manifesting them. While he agrees with the Saiva idealist that

appearances have no independent existence apart from their appearing to consciousness, he maintains that for this reason they are unreal. The creativity of consciousness consists in its diversification in many modes having apparent externality; it is not a creation of objects. While the Kashmiri Saivite agrees that the world is pure conscious- ness alone, he maintains that it is such because it is a real creation of consciousness. The effect is essentially identical with the cause and shares in its reality. Matter and the entire universe are absolutely real, as 'congealed' (sty ana) or 'contracted' (samkucita) forms of consciousness. "This God of consciousness", writes Ksemaraja, "generates the universe and its form is a condensation of His own essence (rasa)" m By boiling sugarcane juice it condenses to form treacle, brown sugar and candy which retains its sweetness. Similarly, consciousness abides unchanged even though it assumes the concrete material form of the five gross elements. The same reality thus abides equally in gross and subtle forms. Consequently no object is totally insentient. Even stones bear a trace (vasana) of consciousness, although it is not clearly apparent because it is not associated with the vital breath (prana) and other components of a psycho-physical organism. Somananda goes so far as to affirm that physical objects, far from being insentient, can only exist insofar as they are aware of themselves as existing. The jar performs its function because it knows itself to be its agent. Indeed, all things are pervaded by consciousness and at one with it and hence share in its omniscience. Thus, Siva, Who perceives Himself in the form of physical objects, is the one ultimate reality. "The jar knows because it is of my nature", writes Somananda, "and I know it because I am of the jar's nature. I know because I am of Sadasiva's nature and He knows because He is of my nature; Yajnadatta [knows] because he is of Siva's nature and Siva [knows] because He is of Yajnadatta's nature". Integral Monism of Kashmiri Saivism Everything in this sense is directly perceived by absolute conscious- ness, and this direct perception (pratyaksa) unifies the knowable into a single, undivided whole. This is the central concept behind a doctrine originally expounded by Narasimha called 'the non-dualism of direct perception' (pratyaksadvaita). This states that consciousness is essentially perceptive and that its perception of all things operates throughout the universe. Insofar

as phenomena are clearly evident (sphufa) to us, everything is directly perceived by absolute consciousness, with which our individual consciousness is identical. This direct perception unfolds everywhere; the one true reality, it is alone and without companion or rival (nihsapatna). Even though it remains one, it can, by its very nature, perceive distinctions (bheda) between one entity and another, without this engendering any division within it. We distinguish between two entities in empirical terms on the basis of their mutual exclusion (anyonyabhava). The relative distinction {bheda) between them is essentially the perceived difference between their respective characteristics. Despite this difference they are united within the purview of a single cognition insofar as they are equally both manifest appearances. This cognition is the undivided essence (rasa) or 'own nature' (svabhava) of both. Encompassed by the 'fire of consciousness', there is no essential difference between them. Just as when an emerald and ruby reflect each other's light, the ruby is reddish-green and the emerald greenish-red, similarly everything is connected with everything else as part of the single variegated (vicitra) cognition of absolute consciousness. Mahesvarananda writes: The Supreme Lord's unique state of emotivity (asadharanabhava) is the outpouring of pure Being (mahasatta). It is manifest as the brilliance (sphuratta) of the universe which, if we ponder deeply, [is realized to be] the single flavour (ekarasa) of the essence of Beauty which is the vibration of the bliss of one's own nature. In this way all things are in reality one although divided from the one another sharing as they do the 'single flavour' (ekarasa) of the pure vibration of consciousness. Kashmiri Saiva Realism Kashmiri Saivism as a whole has been variously called a form of 'realistic idealism', 'monistic idealism', ‘idealistic monism' and 'concrete monism'. It is easy to understand why Kashmiri Saivism is The Doctrine of Vibration said to be 'idealistic' and 'monistic', but in what sense is it also 'realistic*? The answer to this question is of no small importance in trying to under- stand the central idea behind its metaphysics and the fundamental importance of the concept of Spanda, in this seemingly impossible marriage between monistic idealism and pluralistic realism. The Kashmiri Saiva approach understands the world to be a symbol of the absolute, that is, as the manner in which it presents itself to us. Again we can

contrast this view with that of the Advaita Vedanta. The Advaita Vedanta understands the world to be an expression of the absolute insofar as it exists by virtue of the absolute's Being. Being is understood to be the real unity which underlies empirically manifest separateness and as such is never empirically manifest. It is only transcendentally actual as 'being-in-itself. The Kashmiri Saiva position represents, in a sense, a reversal of this point of view. The nature of the absolute, and also that of Being, is onceived as an eternal becoming (satatodita), a dynamic flux or Spanda, 'the agency of the act of being'. It is dentified with the concrete actuality of the fact of appearing, not passive unmanifest Being. Appearance (abhasa) alone is real.ppearing (prakasamanatva) is equivalent to the fact of being (astitva). Ksemaraja writes in his commentary on the Stanzas on Vibration: Indeed, all things are manifest because they are nothing but manifestation. The point being that nothing is manifest apart from manifestation The absolutely unmanifest, from this point of view, can have as little existence as the space in a lattice window of a sky-palace. Nay, even less, because even that space can appear as an imagined mage manifest within consciousness. Everything is real according to the manner in which it appears. Even an illusion is in this sense real, insofar as it appears and is known in the manner in which it appears. The empirical and the real are identical categories of thought. As Abhinava says: Thus this is the supreme doctrine (upanisatf), namely that, when- ever and in whatever form [an entity] appears, that then is its particular nature. Perhaps at this stage a brief comparison with Heidegger's ideas might prove to be enlightening and not altogether out of place. According to Heidegger's phenomenology of Being, reality is intelligible in a two-fold Integral Monism of Kashmiri Saivism manner as 'phenomenon' and 'logos'. Heidegger defines what he means by 'phenomenon' as: "that-which-shows-itself. The manifest . . . phenomena are then the collection of that which lies open in broad daylight or can be brought to the light of day — what the Greeks at times implicitly identified as 'ta onta' (the things-which-are)".

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb) heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna**

**Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version) **Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb) strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel

period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar (=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI: 10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills (SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear

differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black**

**holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-

relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb)

the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this

version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstate thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11 August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of

entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.** Characterizing the behaviour of strongly coupled

quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at λ=λc, the correlation length diverges as ξ∼|λ−λc|−ν, while (e) the gap between the ground state and first excited state vanishes as (e) ΔE∼ξ−z∼|λ−λc|zν. Exactly at the critical point λ=λc, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as ω∼|k−k0|z. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-**

**equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004 STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element.

**Model has been applied for the following results:**

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| **Searches for Sterile Neutrinos with the IceCube Detector M. G. Aartsen et al. (IceCube Collaboration) Phys.**<br><br>**Rev. Lett 117, 071801 – Published 8 August 2016 Physics logo See Viewpoint: Hunting the Sterile Neutrino** |
| **(1) The IceCube neutrino telescope at the South Pole has measured the atmospheric muon neutrino spectrum as (=) a function of zenith angle and (e&eb) energy in the approximate 320 GeV to 20 TeV range, to search for (e) the oscillation signatures of (e) light sterile neutrinos.** |
| **(2) No evidence for anomalous νμ or ν¯μ disappearance is observed in (eb) either of two independently developed analyses, each using one year of (e) atmospheric neutrino data.** |
| **(3) New exclusion limits are placed on (e&eb) the parameter space of (e) the 3+1 model, in which muon antineutrinos experience (e) a strong Mikheyev-Smirnov-Wolfenstein-resonant oscillation.** |
| **(4) The exclusion limits extend to (e) sin22024≤0.02 at Δm2∼0.3 eV2 at the 90% confidence level.** |
| **(5) The allowed region from (e) global analysis of appearance experiments, including (e) LSND and MiniBooNE, is excluded at approximately the 99% confidence level for (e) the global best-fit value of \|Ue4\|2. Received 5 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.117.071801 © 2016 American Physical Society** |
| **Synopsis: Searching for Majorana Neutrinos August 16, 2016The KamLAND-Zen collaboration has run the most sensitive search to date for a radioactive decay that could reveal whether neutrinos are Majorana fermions. Synopsis figureKamLAND-Zen Collaboration** |
| **(6) Are neutrinos and (=) antineutrinos the same particles, as hypothesized by (e) physicist Ettore Majorana?** |
| **(7) If true, this "Majorana" nature of neutrinos could explain (eb) why neutrinos have (e) mass and why the Universe has (e) more matter than antimatter.** |
| **(8) Using an underground neutrino detector in Japan, the KamLAND-Zen collaboration has carried out the most sensitive search to date for a process called neutrinoless double-beta decay, which could confirm (eb) Majorana's hypothesis.** |
| **(9) The results set stringent constraints on the possible properties of (e) Majorana neutrinos.** |
| **(10) Double-beta decay is (=) a nuclear process in which two neutrons turn into (e&eb) two protons (or vice versa), emitting (eb) two antineutrinos.** |
| **(11) If neutrinos are Majorana fermions, (eb) the two antineutrinos could annihilate each other, and some double-beta decays would emit (eb) no antineutrinos.** |
| **(12) However, neutrinoless decay is believed to be (=) extremely rare. Searches of the process using (e) large amounts of the 35 naturally occurring (eb) isotopes that can undergo (e&eb) double-beta decay (such as xenon-136) have so far come up empty.** |

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| **(13) The KamLAND-Zen collaboration boosted their experimental sensitivity to (e) neutrinoless double-beta decays through (e&eb) a combination of factors.** |
| **(14) The researchers deployed clean detectors with (e&eb) very low background noise.** |
| **(15) And they used an unprecedented amount of xenon-136, which they purified to (e) remove any radioactive contaminants that could produce (eb) unwanted signals in the detectors.** |
| **(16) Combining data from several years, they improved (eb) the limit on the probability of neutrinoless double-beta decay by (e) sixfold compared to previous searches.** |
| **(17) This probability corresponds to (e&eb) a decay lifetime whose staggering value exceeds (e) 1026 years.** |
| **(18) This measurement allowed (eb) the researchers to set (eb) the most stringent upper limit on (e&eb) the possible mass of Majorana neutrinos (the particle must be lighter than 61–165 meV). This research is published in Physical Review Letters. –Matteo Rini Matteo Rini is the Deputy Editor of Physics.** |
| **How to Reach a Thousand-Second in-Plane Polarization Lifetime with 0.97−GeV/c Deuterons in a Storage Ring G. Guidoboni et al. (JEDI Collaboration) Phys. Rev. Lett. 117, 054801 – Published 28 July 2016 Physics logo See Synopsis: Deuterons Spin together for a Thousand Seconds** |
| **(19) Authors observe a deuteron beam polarization lifetime near 1000 s in (eb) the horizontal plane of (e) a magnetic storage ring (COSY).** |
| **(20) This long spin coherence time is maintained through (e&eb) a combination of beam bunching, electron cooling, (e&eb) sextupole field corrections, and (e&eb) the suppression of collective effects through (e&eb) beam current limits.** |
| **(21) This record lifetime is required for (e) a storage ring search for (e) an intrinsic electric dipole moment on (e&eb) the deuteron at a statistical sensitivity level approaching 10−29 e cm.** |
| **Synopsis: Deuterons Spin together for a Thousand Seconds July 28, 2016 Researchers set a new record for the in-plane spin-alignment lifetime of deuterons circulating in a magnetic storage ring. Synopsis figure E. Stephenson/Indiana University Center for Spacetime Symmetries Sime sentences and figures are deleted due to spatial restrictions.** |
| **(22) The electric dipole moment (EDM) of (e) a fundamental particle or a simple nucleus is believed to be (=) extremely small.** |
| **(23) One might therefore think that it isn't worth measuring this quantity. But it is. A nonzero EDM would signal (eb) a broken symmetry of (e) nature known as CP violation, which is required to (e) explain why the Universe has (e) more matter than antimatter.** |
| **(24) Edward Stephenson from Indiana University, Bloomington, and colleagues has now made a significant stride towards determining (eb) the EDM of the deuteron using (e) a deuteron beam circulating in (e&eb) a magnetic storage ring.** |

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| **(25) They have been able to preserve the alignment of the deuterons' spins for (e) a chunk of time that may be long enough to determine (eb) the deuteron's EDM.** |
| **(26) In experiments of this kind, the EDM can be determined by (e) measuring the rotation of (e) the initially aligned spins of (e) the beam particles as they travel around (e&eb) the ring.** |
| **(27) To obtain a measurement sufficiently accurate to explore (e&eb) CP violation predicted by (e) theories beyond the standard model, the spin alignment has to be (=) preserved for about 1000 s.** |
| **(28) But changes in particle speed—associated with (e&eb) position or direction variations among (e&eb) the particles in the beam—lead to (eb) a rapid loss of the alignment in the plane of the ring.** |
| **(29) Stephenson and co-workers have managed to retain (eb)the alignment of deuterons for the required 1000 s, which is 3 orders of magnitude longer (eb) than the previous record (obtained for an electron beam).** |
| **(30) Their success resulted from (e) capturing the beam into a tight bunch to reduce (e) position and direction spread and from (e) shaping the ring's magnetic field to(e) ensure all of the deuteron orbits around (e&eb) the ring have the same length, despite (e) direction variations. This research is published in Physical Review Letters. –Ana Lopes Ana Lopes is a Senior Editor of Physics.** |
| **Cotunneling Drag Effect in Coulomb-Coupled Quantum Dots A. J. Keller, J. S. Lim, David Sánchez, Rosa López, S. Amasha, J. A. Katine, Hadas Shtrikman, and D. Goldhaber-Gordon Phys. Rev. Lett. 117, 066602 – Published 4 August 2016 Physics logo See Synopsis: Coulomb Drag in a Double Dot** |
| **(31) In Coulomb drag, a current flowing in one conductor can induce (eb) a voltage across an adjacent conductor via (e&eb) the Coulomb interaction.** |
| **(32) The mechanisms yielding drag effects are not always understood, even though drag effects are (=) sufficiently general to be (=) seen in many low-dimensional systems.** |
| **(33) In this Letter, authors observe Coulomb drag in (eb) a Coulomb-coupled double quantum dot and, through (e&eb) both experimental and theoretical arguments, identify cotunneling as (=) essential to obtaining a correct qualitative understanding of the drag behavior.** |
| **Synopsis: Coulomb Drag in a Double Dot August 4, 2016 Electric current passing through a quantum dot can generate current in a nearby dot through a coordinated tunneling mechanism. Synopsis figureA. J. Keller et al., Phys. Rev. Lett. (2016)** |
| **(34) Electrons moving through (e&eb) a conductor can induce (eb) the movement of electrons of a nearby, but isolated, conductor.** |
| **(35) This so-called Coulomb drag is well understood in 2D layered materials, but the theory for nanowires and quantum dots does not (e) fully explain experiments.** |
| **(36) A new study of Coulomb drag in two nearby quantum dots shows (eb) that simultaneous tunneling of electrons—called cotunneling—can help explain the low-temperature behavior of (e) the system.** |

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| **(37) Electric current flows through (e&eb) a quantum dot by the tunneling of (e) electrons.** |
| **(38) When two dots are placed next to each other, the active, or lead, dot can drag current through (e&eb) the passive dot by (e) a process called sequential tunneling.** |
| **(39) In this case, an electron first tunnels onto (e&eb) the active dot, increasing (eb) the probability of tunneling onto (e&eb) the passive dot.** |
| **(40) Thus the passive dot follows (e&eb) the active dot's lead, but always one step behind.** |
| **(41) Sequential tunneling explains (eb) Coulomb drag observations at (eb) high temperature, but no studies have investigated cold temperatures where (e) quantum coherence effects may play a role.** |
| **(42) Andrew Keller from the California Institute of Technology, Pasadena, and his colleagues explored Coulomb drag at (eb) sub-Kelvin temperatures in a pair of quantum dots coupled by (e) capacitance.** |
| **(43) The team applied a current-driving voltage across (e&eb) one dot and detected a drag current in (eb) the other dot, even in regimes where (e) the driving voltage was very small.** |
| **(44) This conflicts with (e&eb) the sequential tunneling model, which predicts (eb) a threshold driving voltage below which (e) the drag current should be zero.** |
| **(45) The team showed that the inclusion of cotunneling into (eb) the drag theory explains the observations.** |
| **(46) Cotunneling could be important to efforts aimed at harvesting a unidirectional current from (e) a fluctuating voltage supply. This research is published in Physical Review Letters. –Michael Schirber Michael Schirber is a Corresponding Editor for Physics based in Lyon, France.** |
| **Maybe it wasn't the Higgs particle after all November 7, 2014 Maybe it wasn't the Higgs particle after all Associate Professor Mads Toudal Frandsen, University of Southern Denmark Credit: University of Southern Denmark** |
| **(47) Last year CERN announced the finding of a new elementary particle, the Higgs particle. But maybe it wasn't the Higgs particle, maybe it just looks like it. And maybe it is not alone. Many calculations indicate that the particle discovered last year in the CERN particle accelerator was (=) indeed the famous Higgs particle.** |
| **(48) Physicists agree that the CERN experiments did find a new particle that had never been seen before, but according to (e&eb) an international research team, there is (=) no conclusive evidence that the particle was indeed the Higgs particle.The research team has scrutinized the existing scientific data from CERN about the newfound particle and published their analysis in the journal Physical Review D. A member of this team is Mads Toudal Frandsen, associate professor at the Center for Cosmology and Particle Physics Phenomenology, Department of Physics, Chemistry and Pharmacy at the University of Southern Denmark. "The CERN data is generally taken as (=) evidence that the particle is the Higgs particle.** |

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| **(49) It is true that the Higgs particle can explain (eb) the data but there can be other explanations, we would also get this data from other particles", Mads Toudal Frandsen explains.** |
| **(50) The researchers' analysis does not (e) debunk the possibility that CERN has discovered the Higgs particle. That is still possible - but it is equally possible that it is a different kind of particle. "The current data is not precise enough to determine exactly what the particle is. It could be a number of other known particles", says Mads Toudal Frandsen. But if it wasn't the Higgs particle, which was found in CERN's particle accelerator, then what was it? "We believe that it may be a so-called techni-higgs particle. This particle is in some ways similar to the Higgs particle - hence half of the name", says Mads Toudal Frandsen.** |
| **(51) Although the techni-higgs particle and Higgs particle can easily be confused in experiments, they are two very different particles belonging to (e) two very different theories of how the universe was created.** |
| **(52) The Higgs particle is the missing piece in (eb) the theory called the Standard Model.** |
| **(53) This theory describes (eb) three of the four forces of nature. But it does not (e) explain what dark matter is - the substance that makes up (eb) most of the universe. A techni-higgs particle, if it exists, is (=) a completely different thing:** |
| **(54) "A techni-higgs particle is not (e) an elementary particle.** |
| **(55) Instead, it consists of so-called techni-quarks, which we believe are elementary. Techni-quarks may bind together in various ways to form (eb) for instance techni-higgs particles, while other combinations may form (eb) dark matter. Authors therefore expect to find several different particles at the LHC, all built by (e) techni-quarks", says Mads Toudal Frandsen.** |
| **(56) If techni-quarks exist, (eb) there must be a force to bind them together so that (e) they can form particles.** |
| **(57) None of the four known forces of nature (gravity, the electromagnetic force, the weak nuclear force and the strong nuclear force) are (=) any good at binding techni-quarks together.** |
| **(58) There must therefore be (=) a yet undiscovered force of nature. This force is called the the Technicolor force.** |
| **(59) What was found last year in CERN's accelerator could thus be (=) either the Higgs particle of the Standard Model or a light techni-higgs particle, composed of (e) two techni-quarks.** |
| **(60) Mads Toudal Frandsen believes that more data from CERN will probably be able to determine (eb) if it was a Higgs or a techni-higgs particle. If CERN gets an even more powerful accelerator, it will in principle be able to observe techni-quarks directly. More information: Technicolor Higgs boson in the light of LHC data. Phys. Rev. D 90, 035012th Alexander Belyaev, Matthew S. Brown, Roshan Foadi, and Mads T. Frandsen. journals.aps.org/prd/abstract/10.1103/PhysRevD.90.035012. On Arxiv: arxiv.org/abs/1309.2097 Journal reference: Physical Review D search and more info website: Provided by: University of Southern Denmark search and more info website** |

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| **The inflationary brane-antibrane universeClifford P. Burgess1,2, Mahbub Majumdar3, Detlef Nolte1, Fernando Quevedo3, Govindan Rajesh1 and Ren-Jie Zhang1 Published 13 August 2001 • Journal of High Energy Physics, Volume 2001, JHEP07(2001)** |
| **(61) Authors show how the motion through (e&eb) the extra dimensions of a gas of branes and antibranes can, under certain circumstances, produce (eb) an era of inflation as seen by (e) observers trapped on (e&eb) a 3-brane, with (e&eb) the inflaton being the inter-brane separation.** |
| **(62) Although most of discussion refers to (e) arbitrary ps, when we need to be specific they assume that they are (=) D-branes of type-II or type-I string theory.** |
| **(63) For realistic brane couplings, such as those arising in (eb) string theory, the inter-brane potentials are (=) too steep to inflate (eb) the universe for acceptably long times.** |
| **(64) However, for special regions of the parameter space of (e) brane-antibrane positions the brane motion is (=) slow enough for there to be (=) sufficient inflation.** |
| **(65) Inflation would be more generic in models where (e) the inter-brane interactions are much weaker.** |
| **(66) The spectrum of primordial density fluctuations predicted has (e) index n slightly less than 1, and acceptable amplitude, provided (e) that the extra dimensions have (e) linear size 1/r ~ 1012 GeV.** |
| **(67) Reheating occurs as in hybrid inflation, with (e&eb) the tachyonic instability of (e) the brane-antibrane system taking over for (e) small separations.** |
| **(68) The tachyon field can induce (eb) a cascade mechanism within which higher-dimension branes annihilate (e) into lower-dimension ones.** |
| **(69) They argue that such a cascade naturally stops with (e&eb) the production of 3-branes in (eb) 10-dimensional string theory.** |
| **Effect of the neutron lifetime on processes in the early universe A. V. Chechkin, A. V. Ivanchik, A. P. Serebrov, S. V. Bobashev First Online: 27 July 2016 DOI: 10.1134/S1063784216070069 Chechkin, A.V., Ivanchik, A.V., Serebrov, A.P. et al. Tech. Phys. (2016) 61: 1101. Doi: 10.1134/S1063784216070069** |
| **(70) The influence of the neutron lifetime on (e&eb) the abundance of light elements produced during (e&eb) the primordial nucleosynthesis early in (eb) the birth of the Universe is considered.** |
| **(71) Among light elements, namely, D, 3He, 4He, and 7Li, 4He proves to be (=) most sensitive to neutron lifetime τn.** |
| **(72) Astronomic data on the light element abundance also provide (eb) the best accuracy for 4He.** |
| **(73) The solution of a number of problems discussed in this paper requires (e) improving the accuracy of observations for the 4He abundance and refining the value of τn.** |

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| **Application of a Theory and Simulation based Convective Boundary Mixing model for AGB Star Evolution and Nucleosynthesis U. Battino, M. Pignatari, C. Ritter, F. Herwig, P. Denisenkov, J.W. Den Hartogh, R. Trappitsch, R. Hirschi, B. Freytag, F. Thielemann, B. Paxton (Submitted on 19 May 2016)** |
| **(74) The s-process nucleosynthesis in Asymptotic Giant Branch (AGB) stars depends on (e&eb) the modeling of (e) convective boundaries.** |
| **(75) Authors' present models and s-process simulations that adopt (e&eb) a treatment of convective boundaries based on (e) the results of hydrodynamic simulations and on (e&eb) the theory of mixing due to (e) gravity waves in (eb) the vicinity of convective boundaries.** |
| **(76) Hydrodynamics simulations suggest (eb) the presence of convective boundary mixing (CBM) at (eb) the bottom of the thermal pulse-driven convective zone.** |
| **(77) Similarly, convection-induced mixing processes are proposed for (e) the mixing below the convective envelope during (e&eb) third dredge-up where (e) the 13C pocket for (e) the s process in AGB stars forms.** |
| **(78) In this work apply apply a CBM model motivated by (e) simulations and theory to (e&eb) models with initial mass M = 2 and M = 3M?, and with (e&eb) initial metal content Z = 0:01 and Z = 0:02.** |
| **(79) As reported previously, the He-intershell abundance of 12C and 16O are increased by (e) CBM at the bottom of (e) pulse-driven convection zone.** |
| **(80) This mixing is affecting (e&eb) the 22Ne (alpha, n) 25Mg activation and (e&eb) the s-process effciency in the 13C-pocket.** |
| **(81) In our model CBM at the bottom of the convective envelope during (e&eb) the third dredgeup represents (eb) gravity wave mixing.** |
| **(82) Authors take further into account that hydrodynamic simulations indicate (eb) a declining mixing efficiency already about (e&eb) a pressure scale height from (e) the convective boundaries, compared to (e&eb) mixing-length theory.** |
| **(83) The final s-process abundances are characterized by (e) 0:36 < [s=Fe] < 0:78 and the heavy-to-light s-process ratio is (=) 0:23 < [hs=ls] < 0:45.** |
| **(84) Finally, authors compare results with (e&eb) stellar observations, pre-solar grain measurements and previous work. Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI:10.3847/0004-637X/827/1/30 Cite as: arXiv:1605.06159 [astro-ph.SR] (or arXiv:1605.06159v1 [astro-ph.SR] for this version)** |
| **MESA and NuGrid simulations of classical novae: CO and ONe nova nucleosynthesis Pavel A. Denissenkov (UVic), James W. Truran (Chicago), Marco Pignatari (Basel), Reto Trappitsch (Chicago), Christian Ritter (Frankfurt), Falk Herwig (UVic), Umberto Battino (Basel), Kiana Setoodehnia (McMaster), B. Paxton (UC Santa Barbara)** |

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| **(85) Classical novae are the result of (e) thermonuclear flashes of hydrogen accreted by (e) CO or ONe white dwarfs, leading (eb) eventually to the dynamic ejection of (e) the surface layers.** |
| **(86) These are observationally known to be (=) enriched in heavy elements, such as (=) C, O and Ne that must originate in (eb) layers below the H-flash convection zone.** |
| **(87) Building on their previous work, they now present stellar evolution simulations of (e) ONe novae and provide (eb) a comprehensive comparison of models with (e&eb) published ones.** |
| **(88) Some of authors models include exponential convective boundary mixing to account for (e) the observed enrichment of (e) the nova ejecta even when accreted material has (e) a solar abundance distribution.** |
| **(89) Models produce (eb) maximum temperature evolution profiles and nucleosynthesis yields (eb) in good agreement with (e&eb) models that generate (eb) enriched ejecta by assuming (eb) that the accreted material was pre-mixed.** |
| **(90) Authors confirm for (e) ONe novae the result we reported previously, i.e.\ we found that 3He could be produced {\it in situ} in (eb) solar-composition envelopes accreted with (e&eb) slow rates (M˙<10−10M⊙/yr) by (e) cold (TWD<107 K) CO WDs, and that convection was triggered by (eb) 3He burning before (e) the nova outburst in that case.** |
| **(91) In addition, authors now find that the interplay between the 3He production and (e&eb) destruction in the solar-composition envelope accreted with (e&eb) an intermediate rate, e.g.\ M˙=10−10M⊙/yr, by the 1.15M⊙ ONe WD with (e&eb) a relatively high initial central temperature, e.g.\ TWD=15×106 K, leads to (eb) the formation of a thick radiative buffer zone that separates (e&eb) the bottom of the convective envelope from (e) the WD surface. (Abridged) Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI: 10.1093/mnras/stu1000 Cite as: arXiv:1303.6265 [astro-ph.SR] (or arXiv:1303.6265v2 [astro-ph.SR] for this version)** |
| **Hybrid C-O-Ne White Dwarfs as Progenitors of Diverse SNe Ia Pavel Denissenkov (UVic), James Truran (U. Chicago), Falk Herwig (UVic), Sam Jones (UVic), Bill Paxton (UC Santa Barbara), Ken'ichi Nomoto (U. Tokyo), Toshio Suzuki (Nihon U.), Hiroshi Toki (Osaka U.)(Submitted on 6 Nov 2014)** |
| **(92) When carbon is ignited off-center in a CO core of a super-AGB star (eb) its burning in a convective shell tends to propagate to (e&eb) the center.** |
| **(93) Whether the C flame will actually be able to reach (eb) the center depends on (e&eb) the efficiency of (e) extra mixing beneath (e&eb) the C convective shell.** |
| **(94) Whereas thermohaline mixing is too inefficient to interfere with the C-flame propagation, (eb) convective boundary mixing can prevent (e) the C burning from reaching (eb) the center.** |
| **(95) As a result, a C-O-Ne white dwarf (WD) is (=) formed, after the star has lost (e) its envelope.** |
| **(96) Such a "hybrid" WD has (e) a small CO core surrounded by (e) a thick ONe zone.** |

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| **(97) In the proposed 1D stellar evolution computations the hybrid WD is allowed to (eb) accrete C-rich material, as if it were in a close binary system and (e&eb) accreted H-rich material from (e) its companion with (e&eb) a sufficiently high rate at which the accreted H would be processed into (e&eb) He under (e&eb) stationary conditions, assuming (eb) that He could then be transformed into (e&eb) C.** |
| **(98) When the mass of the accreting WD approaches (eb) the Chandrasekhar limit, they find a series of convective Urca shell flashes associated with (e&eb) high abundances of 23Na and 25Mg.** |
| **(99) They are followed by off-center C ignition leading to (eb) convection that occupies (e) almost the entire star.** |
| **(100) To model the Urca processes, authors use the most recent well-resolved data for (e) their reaction and (e&eb) neutrino-energy loss rates.** |
| **(101) Because of the emphasized uncertainty of the convective Urca process in hybrid WD models of (e) SN Ia progenitors, authors consider a number of their potentially possible alternative instances for (e) different mixing assumptions, all of which reach (eb) a phase of explosive C ignition, either off or in (eb) the center.** |
| **(102) Our hybrid SN Ia progenitor models have (e) much lower C to O abundance ratios at (eb) the moment of the explosive C ignition than (e) their pure CO counterparts, which may explain (eb) the observed diversity of the SNe Ia. subjects: Solar and Stellar Astrophysics (astro-ph.SR)cite as: arXiv:1411.1471 [astro-ph.SR](or arXiv:1411.1471v1 [astro-ph.SR] for this version)** |
| **Rejecting proposed dense-matter equations of state with quiescent low-mass X-ray binaries Sebastien Guillot, Robert E. Rutledge (Submitted on 15 Sep 2014 (v1), last revised 20 Oct 2014 (this version, v2))** |
| **(103) Neutrons stars are unique laboratories to (e) discriminate between the various proposed equations of (e) state of matter at and above nuclear density.** |
| **(104) One sub-class of neutron stars - those inside (eb) quiescent low-mass X-ray binaries (qLMXBs) - produce (eb) a thermal surface emission from (e) which the neutron star radius (R_NS) can be measured, using (e) the widely accepted observational scenario for qLMXBs, assuming (eb) unmagnetized H atmospheres.** |
| **(105) In a combined spectral analysis, this work first reproduces (eb) a previously published measurement of the \rns, assumed to be (=) the same for all neutron stars, using (e) a slightly expanded data set.** |
| **(106) The radius measured is R_NS = 9.4 +/-1.2 km. On the basis of spectral analysis alone, this measured value is not affected by imposing (e&eb) an assumption of causality in the core.** |
| **(107) However, the assumptions underlying this R_NS measurement would be falsified by (e) the observation of any neutron star with (e&eb) a mass >2.6 Msun, since radii <11 km would be rejected if (e) causality is assumed, which would exclude (eb) most of the R_NS parameter space obtained in this analysis.** |

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| **(108) Finally, this work directly tests (e&eb) a selection of dense matter equations of states: WFF1, AP4, MPA1, PAL1, MS0, and three versions of equations of state produced through (e&eb) chiral effective theory.** |
| **(109) Two of those, MSO and PAL1, are rejected at (eb) the 99% confidence level, accounting for (e) all quantifiable uncertainties, while (e) the other cannot be excluded at >99% certainty. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE) DOI: 10.1088/2041-8205/796/1/L3 Cite as: arXiv:1409.4306 [astro-ph.HE] (or arXiv:1409.4306v2 [astro-ph.HE] for this version** |

## SECTION FOUR

## Molecular Collisions Constitute The Dominant Noise Source

**INTRODUCTION—VARIABLES USED**

**Viewpoint: Measuring Quantum Kicks from a Beam of Light Amir Safavi-Naeini, 348 Via Pueblo Mall, Stanford University, Stanford, CA 94305-4088, USA June 13, 2016• Physics 9, 64 Force sensors levitated by light have reached the quantum regime, in which their sensitivity is limited by the momentum kicks of individual photons. Figure caption expand figure APS/Alan Stonebraker Figures and some sentences are deleted due to spatial restrictions Please excuse me**

(1) This final temperature has (e) a linear dependence on pressure from 10−610−6 to (e) 10−210−2 mbar, the range of trap pressures in (eb) which molecular collisions constitute (e) the dominant noise source.
(2) However, at pressures below 10−610−6 mbar the equilibrium temperature ceases to (e) depend on the pressure.
(3) This is the regime in which photon recoil is (=) the dominant source of heating and force noise.
(4) To measure the photon-recoil heating rate, the ETH group turned off their mechanism for (e) cooling the particle by (e) holding the laser intensity constant and allowing (eb) the particle to heat up through (e&eb) its interaction with (e&eb) the (room-temperature) environment.
(5) They then measured the rate at which this warm-up occurred, over many runs and at ultralow pressures, and determined that its value was on (e&eb) the order of a few tens of kilohertz—an impressively slow rate for (e) a room-temperature experiment.
(6) In addition, the measured values of the photon-recoil heating rate for various experimental parameters agree with (=) quantum predictions.
(7) Since the first proposals for using (e) trapped nanoparticles in (eb) quantum optical and optomechanical experiments [9, 10], major efforts have been directed at using (e) these systems to observe (e) quantum phenomena.
(8) Such sensors might allow us to someday test quantum mechanics itself in the limit of (e) “massive” systems, which, like the nanosphere, consist of (e) hundreds of millions of atoms [11].
(9) By demonstrating a measurement of photon recoil heating, Novotny’s group has unambiguously observed (e) the first signatures of quantum phenomena on (e&eb) optically levitated nanoparticles

**NOTATION**

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| **Module One**<br><br>This final temperature has (e) a linear dependence on pressure from 10−610−6 to (e) 10−210−2 mbar, the range of trap pressures in (eb) which molecular collisions constitute (e) the dominant noise source | |
| $G_{13}$ : Category one of linear dependence on pressure from 10−610−6 to (e) 10−210−2 mbar, the range of trap pressures in (eb) which molecular collisions constitute (e) the dominant noise source<br><br>$G_{14}$ : Category two of SAS<br><br>$G_{15}$ : Category three of SAS | |
| $T_{13}$ : Category one of final temperature<br><br>$T_{14}$ : Category two of SAS<br><br>$T_{15}$ : Category three of SAS | |
| **Module Two**<br><br>This final temperature has a linear dependence on pressure from 10−610−6 to (e) 10−210−2 mbar, the range of trap pressures in (eb) which molecular collisions constitute (e) the dominant noise source | |
| $G_{16}$ : Category one of final temperature has a linear dependence on pressure from 10−610−6 to (e) 10−210−2 mbar, the range of trap pressures<br><br>$G_{17}$ : Category two of SAS<br><br>$G_{18}$ : Category three of SAS | |
| $T_{16}$ : Category one of molecular collisions constitute (e) the dominant noise source<br><br>$T_{17}$ : Category two of SAS<br><br>$T_{18}$ : Category three of SAS | |
| **Module three**<br><br>This final temperature has a linear dependence on pressure from 10−610−6 to (e) 10−210−2 mbar, the range of trap pressures in which molecular collisions constitute (e) the dominant noise source | |
| $G_{20}$ : Category one of dominant noise source<br><br>$G_{21}$ : Category two of SAS<br><br>$G_{22}$ : Category three of SAS | |
| $T_{20}$ : Category one of final temperature has a linear dependence on pressure from 10−610−6 to (e) 10−210−2 mbar, the range of trap pressures in which molecular collisions<br><br>$T_{21}$ : Category two of SAS | |

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| $T_{22}$ : Category three of SAS | |
| **Module four**<br>However, at pressures below 10−610−6 mbar the equilibrium temperature ceases to (e) depend on the pressure | |
| $G_{24}$ : Category one of pressures below 10−610−6 mbar the equilibrium temperature; depend on the pressure<br>$G_{25}$ : Category two of SAS<br>$G_{26}$ : Category three of SAS | |
| $T_{24}$ : Category one of depend on the pressure; pressures below 10−610−6 mbar the equilibrium temperature<br>$T_{25}$ : Category two of SAS<br>$T_{26}$ : Category three of SAS | |
| **Module five**<br>This is the regime in which photon recoil is (=) the dominant source of heating and force noise | |
| $G_{28}$ : Category one of regime in which photon recoil<br>$G_{29}$ : Category two of SAS<br>$G_{30}$ : Category three of SAS | |
| $T_{28}$ : Category one of dominant source of heating and force noise<br>$T_{29}$ : Category two of SAS<br>$T_{30}$ : Category three of SAS | |
| **Module six**<br>To measure the photon-recoil heating rate, the ETH group turned off their mechanism for (e) cooling the particle by (e) holding the laser intensity constant and allowing (eb) the particle to heat up through (e&eb) its interaction with (e&eb) the (room-temperature) environment | |
| $G_{32}$ : Category one of cooling the particle by (e) holding the laser intensity constant and allowing (eb) the particle to heat up through (e&eb) its interaction with (e&eb) the (room-temperature) environment<br>$G_{33}$ : Category two of SAS<br>$G_{34}$ : Category three of SAS | |
| $T_{32}$ : Category one of measure the photon-recoil heating rate, the ETH group turned off their mechanism<br>$T_{33}$ : Category two of SAS | |

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| $T_{34}$ : Category three of SAS | |
| **Module seven**<br><br>To measure the photon-recoil heating rate, the ETH group turned off their mechanism for cooling the particle by (e) holding the laser intensity constant and allowing (eb) the particle to heat up through (e&eb) its interaction with (e&eb) the (room-temperature) environment | |
| $G_{36}$ : Category one of holding the laser intensity constant and allowing (eb) the particle to heat up through (e&eb) its interaction with (e&eb) the (room-temperature) environment<br><br>$G_{37}$ : Category two of SAS<br><br>$G_{38}$ : Category three of SAS | |
| $T_{36}$ : Category one of measure the photon-recoil heating rate, the ETH group turned off their mechanism for cooling the particle<br><br>$T_{37}$ : Category two of SAS<br><br>$T_{38}$ : Category three of SAS | |
| **Module eight**<br><br>To measure the photon-recoil heating rate, the ETH group turned off their mechanism for cooling the particle by holding the laser intensity constant and allowing (eb) the particle to heat up through (e&eb) its interaction with (e&eb) the (room-temperature) environment | |
| $G_{40}$ : Category one of measure the photon-recoil heating rate, the ETH group turned off their mechanism for cooling the particle by holding the laser intensity constant<br><br>$G_{41}$ : Category two of SAS<br><br>$G_{42}$ : Category three of SAS | |
| $T_{40}$ : Category one of particle to heat up through (e&eb) its interaction with (e&eb) the (room-temperature) environment<br><br>$T_{41}$ : Category two of SAS<br><br>$T_{42}$ : Category three of SAS | |
| **Module Nine**<br><br>To measure the photon-recoil heating rate, the ETH group turned off their mechanism for cooling the particle by holding the laser intensity constant and allowing the particle to heat up through (e&eb) its interaction with (e&eb) the (room-temperature) environment | |
| $G_{44}$ : Category one of **measure the photon-recoil heating rate, the ETH group turned off their mechanism for cooling the particle by holding the laser intensity constant and allowing the particle to** | |

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| **heat up;** its interaction with (e&eb) the (room-temperature) environment<br><br>$G_{45}$ : Category two of SAS<br><br>$G_{46}$ : Category three of SAS | |
| $T_{44}$ : Category one of its interaction with (e&eb) the (room-temperature) environment **;measure the photon-recoil heating rate, the ETH group turned off their mechanism for cooling the particle by holding the laser intensity constant and allowing the particle to heat up**<br><br>$T_{45}$ : Category two of SAS<br><br>$T_{46}$ : Category three of SAS | |
| **The Coefficients:** | |
| $(a_{13})^{(1)}, (a_{14})^{(1)}, (a_{15})^{(1)}, (b_{13})^{(1)}, (b_{14})^{(1)}, (b_{15})^{(1)}$ $(a_{16})^{(2)}, (a_{17})^{(2)}, (a_{18})^{(2)}$ $(b_{16})^{(2)}, (b_{17})^{(2)}, (b_{18})^{(2)}$:<br>$(a_{20})^{(3)}, (a_{21})^{(3)}, (a_{22})^{(3)}$ , $(b_{20})^{(3)}, (b_{21})^{(3)}, (b_{22})^{(3)}$<br>$(a_{24})^{(4)}, (a_{25})^{(4)}, (a_{26})^{(4)}, (b_{24})^{(4)}, (b_{25})^{(4)}, (b_{26})^{(4)}, (b_{28})^{(5)}, (b_{29})^{(5)}, (b_{30})^{(5)}, (a_{28})^{(5)}, (a_{29})^{(5)}, (a_{30})^{(5)}$,<br>$(a_{32})^{(6)}, (a_{33})^{(6)}, (a_{34})^{(6)}, (b_{32})^{(6)}, (b_{33})^{(6)}, (b_{34})^{(6)}$<br>$(a_{36})^{(7)}, (a_{37})^{(7)}, (a_{38})^{(7)}, (b_{36})^{(7)}, (b_{37})^{(7)}, (b_{38})^{(7)}$<br>$(a_{40})^{(8)}, (a_{41})^{(8)}, (a_{42})^{(8)}, (b_{40})^{(8)}, (b_{41})^{(8)}, (b_{42})^{(8)}$<br>$(a_{44})^{(9)}, (a_{45})^{(9)}, (a_{46})^{(9)}, (b_{44})^{(9)}, (b_{45})^{(9)}, (b_{46})^{(9)}$<br><br>are Accentuation coefficients<br>$(a_{13}^{'})^{(1)}, (a_{14}^{'})^{(1)}, (a_{15}^{'})^{(1)}, (b_{13}^{'})^{(1)}, (b_{14}^{'})^{(1)}, (b_{15}^{'})^{(1)}, (a_{16}^{'})^{(2)}, (a_{17}^{'})^{(2)}, (a_{18}^{'})^{(2)}$, $(b_{16}^{'})^{(2)}, (b_{17}^{'})^{(2)}, (b_{18}^{'})^{(2)}$<br>$, (a_{20}^{'})^{(3)}, (a_{21}^{'})^{(3)}, (a_{22}^{'})^{(3)}, (b_{20}^{'})^{(3)}, (b_{21}^{'})^{(3)}, (b_{22}^{'})^{(3)}$<br>$(a_{24}^{'})^{(4)}, (a_{25}^{'})^{(4)}, (a_{26}^{'})^{(4)}, (b_{24}^{'})^{(4)}, (b_{25}^{'})^{(4)}, (b_{26}^{'})^{(4)}, (b_{28}^{'})^{(5)}, (b_{29}^{'})^{(5)}, (b_{30}^{'})^{(5)}$ $(a_{28}^{'})^{(5)}, (a_{29}^{'})^{(5)}, (a_{30}^{'})^{(5)}$ ,<br>$(a_{32}^{'})^{(6)}, (a_{33}^{'})^{(6)}, (a_{34}^{'})^{(6)}, (b_{32}^{'})^{(6)}, (b_{33}^{'})^{(6)}, (b_{34}^{'})^{(6)}$<br>$(a_{36}^{'})^{(7)}, (a_{37}^{'})^{(7)}, (a_{38}^{'})^{(7)}, (b_{36}^{'})^{(7)}, (b_{37}^{'})^{(7)}, (b_{38}^{'})^{(7)}$,<br>$(a_{40}^{'})^{(8)}, (a_{41}^{'})^{(8)}, (a_{42}^{'})^{(8)}, (b_{40}^{'})^{(8)}, (b_{41}^{'})^{(8)}, (b_{42}^{'})^{(8)}$,<br>$(a_{44}^{'})^{(9)}, (a_{45}^{'})^{(9)}, (a_{46}^{'})^{(9)}, (b_{44}^{'})^{(9)}, (b_{45}^{'})^{(9)}, (b_{46}^{'})^{(9)}$,<br>are Dissipation coefficients | |
| **Module Numbered One**<br><br>**The differential system of this model is now (Module Numbered one)** | |
| $\frac{dG_{13}}{dt} = (a_{13})^{(1)} G_{14} - \left[(a_{13}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14}, t)\right] G_{13}$ | 1 |
| $\frac{dG_{14}}{dt} = (a_{14})^{(1)} G_{13} - \left[(a_{14}^{'})^{(1)} + (a_{14}^{''})^{(1)}(T_{14}, t)\right] G_{14}$ | 2 |
| $\frac{dG_{15}}{dt} = (a_{15})^{(1)} G_{14} - \left[(a_{15}^{'})^{(1)} + (a_{15}^{''})^{(1)}(T_{14}, t)\right] G_{15}$ | 3 |
| $\frac{dT_{13}}{dt} = (b_{13})^{(1)} T_{14} - \left[(b_{13}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G, t)\right] T_{13}$ | 4 |
| $\frac{dT_{14}}{dt} = (b_{14})^{(1)} T_{13} - \left[(b_{14}^{'})^{(1)} - (b_{14}^{''})^{(1)}(G, t)\right] T_{14}$ | 5 |
| $\frac{dT_{15}}{dt} = (b_{15})^{(1)} T_{14} - \left[(b_{15}^{'})^{(1)} - (b_{15}^{''})^{(1)}(G, t)\right] T_{15}$ | 6 |
| $+(a_{13}^{''})^{(1)}(T_{14}, t) =$ First augmentation factor | |

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| $-(b_{13}^{''})^{(1)}(G,t) =$ First detritions factor | |
| **Module Numbered Two**<br><br>**The differential system of this model is now ( Module numbered two)** | |
| $\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \left[(a_{16}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17},t)\right]G_{16}$ | 7 |
| $\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \left[(a_{17}^{'})^{(2)} + (a_{17}^{''})^{(2)}(T_{17},t)\right]G_{17}$ | 8 |
| $\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \left[(a_{18}^{'})^{(2)} + (a_{18}^{''})^{(2)}(T_{17},t)\right]G_{18}$ | 9 |
| $\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \left[(b_{16}^{'})^{(2)} - (b_{16}^{''})^{(2)}((G_{19}),t)\right]T_{16}$ | 10 |
| $\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \left[(b_{17}^{'})^{(2)} - (b_{17}^{''})^{(2)}((G_{19}),t)\right]T_{17}$ | 11 |
| $\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \left[(b_{18}^{'})^{(2)} - (b_{18}^{''})^{(2)}((G_{19}),t)\right]T_{18}$ | 12 |
| $+(a_{16}^{''})^{(2)}(T_{17},t) =$ First augmentation factor | |
| $-(b_{16}^{''})^{(2)}((G_{19}),t) =$ First detritions factor | |
| **Module Numbered Three**<br><br>**The differential system of this model is now (Module numbered three)** | |
| $\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \left[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21},t)\right]G_{20}$ | 13 |
| $\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \left[(a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)}(T_{21},t)\right]G_{21}$ | 14 |
| $\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \left[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21},t)\right]G_{22}$ | 15 |
| $\frac{dT_{20}}{dt} = (b_{20})^{(3)}T_{21} - \left[(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23},t)\right]T_{20}$ | 16 |
| $\frac{dT_{21}}{dt} = (b_{21})^{(3)}T_{20} - \left[(b_{21}^{'})^{(3)} - (b_{21}^{''})^{(3)}(G_{23},t)\right]T_{21}$ | 17 |
| $\frac{dT_{22}}{dt} = (b_{22})^{(3)}T_{21} - \left[(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23},t)\right]T_{22}$ | 18 |
| $+(a_{20}^{''})^{(3)}(T_{21},t) =$ First augmentation factor | |
| $-(b_{20}^{''})^{(3)}(G_{23},t) =$ First detritions factor | |
| **Module Numbered Four**<br><br>**The differential system of this model is now (Module numbered Four)** | |
| $\frac{dG_{24}}{dt} = (a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25},t)\right]G_{24}$ | 19 |
| $\frac{dG_{25}}{dt} = (a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25},t)\right]G_{25}$ | 20 |
| $\frac{dG_{26}}{dt} = (a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25},t)\right]G_{26}$ | 21 |
| $\frac{dT_{24}}{dt} = (b_{24})^{(4)}T_{25} - \left[(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}((G_{27}),t)\right]T_{24}$ | 22 |
| $\frac{dT_{25}}{dt} = (b_{25})^{(4)}T_{24} - \left[(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}((G_{27}),t)\right]T_{25}$ | 23 |
| $\frac{dT_{26}}{dt} = (b_{26})^{(4)}T_{25} - \left[(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}((G_{27}),t)\right]T_{26}$ | 24 |
| $+(a_{24}^{''})^{(4)}(T_{25},t) =$ **First augmentation factor** | |
| $-(b_{24}^{''})^{(4)}((G_{27}),t) =$ **First detritions factor** | |
| **Module Numbered Five**:<br><br>**The differential system of this model is now (Module number five)** | |

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| $\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}, t)\right]G_{28}$ | 25 |
| $\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29}, t)\right]G_{29}$ | 26 |
| $\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \left[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29}, t)\right]G_{30}$ | 27 |
| $\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \left[(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}((G_{31}), t)\right]T_{28}$ | 28 |
| $\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \left[(b_{29}^{'})^{(5)} - (b_{29}^{''})^{(5)}((G_{31}), t)\right]T_{29}$ | 29 |
| $\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \left[(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}((G_{31}), t)\right]T_{30}$ | 30 |
| $+(a_{28}^{''})^{(5)}(T_{29}, t) =$ **First augmentation factor** | |
| $-(b_{28}^{''})^{(5)}((G_{31}), t) =$ **First detritions factor** | |
| **Module Numbered Six**<br><br>**The differential system of this model is now (Module numbered Six)** | |
| $\frac{dG_{32}}{dt} = (a_{32})^{(6)}G_{33} - \left[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33}, t)\right]G_{32}$ | 31 |
| $\frac{dG_{33}}{dt} = (a_{33})^{(6)}G_{32} - \left[(a_{33}^{'})^{(6)} + (a_{33}^{''})^{(6)}(T_{33}, t)\right]G_{33}$ | 32 |
| $\frac{dG_{34}}{dt} = (a_{34})^{(6)}G_{33} - \left[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33}, t)\right]G_{34}$ | 33 |
| $\frac{dT_{32}}{dt} = (b_{32})^{(6)}T_{33} - \left[(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}((G_{35}), t)\right]T_{32}$ | 34 |
| $\frac{dT_{33}}{dt} = (b_{33})^{(6)}T_{32} - \left[(b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}((G_{35}), t)\right]T_{33}$ | 35 |
| $\frac{dT_{34}}{dt} = (b_{34})^{(6)}T_{33} - \left[(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}((G_{35}), t)\right]T_{34}$ | 36 |
| $+(a_{32}^{''})^{(6)}(T_{33}, t) =$ **First augmentation factor** | |
| **Module Numbered Seven:**<br><br>**The differential system of this model is now (Seventh Module)** | |
| $\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \left[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37}, t)\right]G_{36}$ | 37 |
| $\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \left[(a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37}, t)\right]G_{37}$ | *38* |
| $\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \left[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37}, t)\right]G_{38}$ | *39* |
| $\frac{dT_{36}}{dt} = (b_{36})^{(7)}T_{37} - \left[(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}((G_{39}), t)\right]T_{36}$ | *40* |
| $\frac{dT_{37}}{dt} = (b_{37})^{(7)}T_{36} - \left[(b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}((G_{39}), t)\right]T_{37}$ | *41* |
| $\frac{dT_{38}}{dt} = (b_{38})^{(7)}T_{37} - \left[(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}((G_{39}), t)\right]T_{38}$ | *42* |
| $+(a_{36}^{''})^{(7)}(T_{37}, t) =$ **First augmentation factor** | |
| **Module Numbered Eight**<br><br>The differential system of this model is now | |
| $\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \left[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41}, t)\right]G_{40}$ | *43* |
| $\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \left[(a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41}, t)\right]G_{41}$ | *44* |
| $\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \left[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41}, t)\right]G_{42}$ | *45* |
| $\frac{dT_{40}}{dt} = (b_{40})^{(8)}T_{41} - \left[(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}((G_{43}), t)\right]T_{40}$ | *46* |
| $\frac{dT_{41}}{dt} = (b_{41})^{(8)}T_{40} - \left[(b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}((G_{43}), t)\right]T_{41}$ | *47* |

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| $\frac{dT_{42}}{dt} = (b_{42})^{(8)}T_{41} - \left[(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}((G_{43}),t)\right]T_{42}$ | 48 |
| **Module Numbered Nine**<br><br>The differential system of this model is now | |
| $\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \left[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45},t)\right]G_{44}$ | 49 |
| $\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \left[(a_{45}^{'})^{(9)} + (a_{45}^{''})^{(9)}(T_{45},t)\right]G_{45}$ | 50 |
| $\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \left[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45},t)\right]G_{46}$ | 51 |
| $\frac{dT_{44}}{dt} = (b_{44})^{(9)}T_{45} - \left[(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}((G_{47}),t)\right]T_{44}$ | 52 |
| $\frac{dT_{45}}{dt} = (b_{45})^{(9)}T_{44} - \left[(b_{45}^{'})^{(9)} - (b_{45}^{''})^{(9)}((G_{47}),t)\right]T_{45}$ | 53 |
| $\frac{dT_{46}}{dt} = (b_{46})^{(9)}T_{45} - \left[(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}((G_{47}),t)\right]T_{46}$ | 54 |
| $+(a_{44}^{''})^{(9)}(T_{45},t) =$ **First augmentation factor** | |
| $-(b_{44}^{''})^{(9)}((G_{47}),t) =$ **First detrition factor** | |
| $\frac{dG_{13}}{dt} = (a_{13})^{(1)}G_{14} - \left[\begin{array}{ccc} (a_{13}^{'})^{(1)}\boxed{+(a_{13}^{''})^{(1)}(T_{14},t)} & \boxed{+(a_{16}^{''})^{(2,2,)}(T_{17},t)} & \boxed{+(a_{20}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{24}^{''})^{(4,4,4,4,)}(T_{25},t)} & \boxed{+(a_{28}^{''})^{(5,5,5,5,)}(T_{29},t)} & \boxed{+(a_{32}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{36}^{''})^{(7,7)}(T_{37},t)} & \boxed{+(a_{40}^{''})^{(8,8)}(T_{41},t)} & \boxed{+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{array}\right]G_{13}$ | 55 |
| $\frac{dG_{14}}{dt} = (a_{14})^{(1)}G_{13} - \left[\begin{array}{ccc} (a_{14}^{'})^{(1)}\boxed{+(a_{14}^{''})^{(1)}(T_{14},t)} & \boxed{+(a_{17}^{''})^{(2,2,)}(T_{17},t)} & \boxed{+(a_{21}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{25}^{''})^{(4,4,4,4,)}(T_{25},t)} & \boxed{+(a_{29}^{''})^{(5,5,5,5,)}(T_{29},t)} & \boxed{+(a_{33}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{37}^{''})^{(7,7)}(T_{37},t)} & \boxed{+(a_{41}^{''})^{(8,8)}(T_{41},t)} & \boxed{+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{array}\right]G_{14}$ | 56 |
| $\frac{dG_{15}}{dt} = (a_{15})^{(1)}G_{14} - \left[\begin{array}{ccc} (a_{15}^{'})^{(1)}\boxed{+(a_{15}^{''})^{(1)}(T_{14},t)} & \boxed{+(a_{18}^{''})^{(2,2,)}(T_{17},t)} & \boxed{+(a_{22}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{26}^{''})^{(4,4,4,4,)}(T_{25},t)} & \boxed{+(a_{30}^{''})^{(5,5,5,5,)}(T_{29},t)} & \boxed{+(a_{34}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{38}^{''})^{(7,7)}(T_{37},t)} & \boxed{+(a_{42}^{''})^{(8,8)}(T_{41},t)} & \boxed{+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{array}\right]G_{15}$ | 57 |
| Where $\boxed{(a_{13}^{''})^{(1)}(T_{14},t)}$, $\boxed{(a_{14}^{''})^{(1)}(T_{14},t)}$, $\boxed{(a_{15}^{''})^{(1)}(T_{14},t)}$ are first augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{16}^{''})^{(2,2,)}(T_{17},t)}$, $\boxed{+(a_{17}^{''})^{(2,2,)}(T_{17},t)}$, $\boxed{+(a_{18}^{''})^{(2,2,)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{20}^{''})^{(3,3,)}(T_{21},t)}$, $\boxed{+(a_{21}^{''})^{(3,3,)}(T_{21},t)}$, $\boxed{+(a_{22}^{''})^{(3,3,)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{24}^{''})^{(4,4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{25}^{''})^{(4,4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{26}^{''})^{(4,4,4,4,)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{28}^{''})^{(5,5,5,5,)}(T_{29},t)}$, $\boxed{+(a_{29}^{''})^{(5,5,5,5,)}(T_{29},t)}$, $\boxed{+(a_{30}^{''})^{(5,5,5,5,)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{32}^{''})^{(6,6,6,6,)}(T_{33},t)}$, $\boxed{+(a_{33}^{''})^{(6,6,6,6,)}(T_{33},t)}$, $\boxed{+(a_{34}^{''})^{(6,6,6,6,)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{38}^{''})^{(7,7)}(T_{37},t)}$ $\boxed{+(a_{37}^{''})^{(7,7)}(T_{37},t)}$ $\boxed{+(a_{36}^{''})^{(7,7)}(T_{37},t)}$ are seventh augmentation coefficient for 1,2,3<br>$\boxed{+(a_{40}^{''})^{(8,8)}(T_{41},t)}$ $\boxed{+(a_{41}^{''})^{(8,8)}(T_{41},t)}$ $\boxed{+(a_{42}^{''})^{(8,8)}(T_{41},t)}$ are eight augmentation coefficient for 1,2,3 | |

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| $+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficient for 1,2,3 | |
| $\frac{dT_{13}}{dt} = (b_{13})^{(1)}T_{14} - \begin{bmatrix} (b_{13}^{'})^{(1)}\left[-(b_{13}^{''})^{(1)}(G,t)\right] & \left[-(b_{16}^{''})^{(2,2,)}(G_{19},t)\right] & \left[-(b_{20}^{''})^{(3,3,)}(G_{23},t)\right] \\ \left[-(b_{24}^{''})^{(4,4,4,4,)}(G_{27},t)\right] & \left[-(b_{28}^{''})^{(5,5,5,5,)}(G_{31},t)\right] & \left[-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)\right] \\ \left[-(b_{36}^{''})^{(7,7,)}(G_{39},t)\right] & \left[-(b_{40}^{''})^{(8,8)}(G_{43},t)\right] & \left[-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{13}$ | 58 |
| $\frac{dT_{14}}{dt} = (b_{14})^{(1)}T_{13} - \begin{bmatrix} (b_{14}^{'})^{(1)}\left[-(b_{14}^{''})^{(1)}(G,t)\right] & \left[-(b_{17}^{''})^{(2,2,)}(G_{19},t)\right] & \left[-(b_{21}^{''})^{(3,3,)}(G_{23},t)\right] \\ \left[-(b_{25}^{''})^{(4,4,4,4,)}(G_{27},t)\right] & \left[-(b_{29}^{''})^{(5,5,5,5,)}(G_{31},t)\right] & \left[-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)\right] \\ \left[-(b_{37}^{''})^{(7,7,)}(G_{39},t)\right] & \left[-(b_{41}^{''})^{(8,8)}(G_{43},t)\right] & \left[-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{14}$ | 59 |
| $\frac{dT_{15}}{dt} = (b_{15})^{(1)}T_{14} - \begin{bmatrix} (b_{15}^{'})^{(1)}\left[-(b_{15}^{''})^{(1)}(G,t)\right] & \left[-(b_{18}^{''})^{(2,2,)}(G_{19},t)\right] & \left[-(b_{22}^{''})^{(3,3,)}(G_{23},t)\right] \\ \left[-(b_{26}^{''})^{(4,4,4,4,)}(G_{27},t)\right] & \left[-(b_{30}^{''})^{(5,5,5,5,)}(G_{31},t)\right] & \left[-(b_{34}^{''})^{(6,6,6,6,)}(G_{35},t)\right] \\ \left[-(b_{38}^{''})^{(7,7,)}(G_{39},t)\right] & \left[-(b_{42}^{''})^{(8,8)}(G_{43},t)\right] & \left[-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{15}$ | *60* |
| Where $-(b_{13}^{''})^{(1)}(G,t)$, $-(b_{14}^{''})^{(1)}(G,t)$, $-(b_{15}^{''})^{(1)}(G,t)$ are first detrition coefficients for category 1, 2 and 3<br>$-(b_{16}^{''})^{(2,2,)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3<br>$-(b_{20}^{''})^{(3,3,)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3<br>$-(b_{24}^{''})^{(4,4,4,4,)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3<br>$-(b_{28}^{''})^{(5,5,5,5,)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3<br>$-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,)}(G_{35},t)$ are sixth detrition coefficients for category 1, 2 and 3<br>$-(b_{37}^{''})^{(7,7,)}(G_{39},t)$, $-(b_{36}^{''})^{(7,7,)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3<br>$-(b_{40}^{''})^{(8,8)}(G_{43},t)$ $-(b_{41}^{''})^{(8,8)}(G_{43},t)$ $-(b_{42}^{''})^{(8,8)}(G_{43},t)$ are eight detrition coefficients for category 1, 2 and 3<br>$-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3 | |
| | |
| $\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \begin{bmatrix} (a_{16}^{'})^{(2)}\left[+(a_{16}^{''})^{(2)}(T_{17},t)\right] & \left[+(a_{13}^{''})^{(1,1,)}(T_{14},t)\right] & \left[+(a_{20}^{''})^{(3,3,3)}(T_{21},t)\right] \\ \left[+(a_{24}^{''})^{(4,4,4,4,4)}(T_{25},t)\right] & \left[+(a_{28}^{''})^{(5,5,5,5,5)}(T_{29},t)\right] & \left[+(a_{32}^{''})^{(6,6,6,6,6)}(T_{33},t)\right] \\ \left[+(a_{36}^{''})^{(7,7,7)}(T_{37},t)\right] & \left[+(a_{40}^{''})^{(8,8,8)}(T_{41},t)\right] & \left[+(a_{44}^{''})^{(9,9)}(T_{45},t)\right] \end{bmatrix} G_{16}$ | *61* |

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| $\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \left[\begin{array}{ccc} (a_{17}^{'})^{(2)}\boxed{+(a_{17}^{''})^{(2)}(T_{17},t)} & \boxed{+(a_{14}^{''})^{(1,1,)}(T_{14},t)} & \boxed{+(a_{21}^{''})^{(3,3,3)}(T_{21},t)} \\ \boxed{+(a_{25}^{''})^{(4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{29}^{''})^{(5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{33}^{''})^{(6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{37}^{''})^{(7,7,7)}(T_{37},t)} & \boxed{+(a_{41}^{''})^{(8,8,8)}(T_{41},t)} & \boxed{+(a_{45}^{''})^{(9,9)}(T_{45},t)} \end{array}\right] G_{17}$ | 62 |
| $\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \left[\begin{array}{ccc} (a_{18}^{'})^{(2)}\boxed{+(a_{18}^{''})^{(2)}(T_{17},t)} & \boxed{+(a_{15}^{''})^{(1,1,)}(T_{14},t)} & \boxed{+(a_{22}^{''})^{(3,3,3)}(T_{21},t)} \\ \boxed{+(a_{26}^{''})^{(4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{30}^{''})^{(5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{34}^{''})^{(6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{38}^{''})^{(7,7,7)}(T_{37},t)} & \boxed{+(a_{42}^{''})^{(8,8,8)}(T_{41},t)} & \boxed{+(a_{46}^{''})^{(9,9)}(T_{45},t)} \end{array}\right] G_{18}$ | 63 |
| Where $\boxed{+(a_{16}^{''})^{(2)}(T_{17},t)}$, $\boxed{+(a_{17}^{''})^{(2)}(T_{17},t)}$, $\boxed{+(a_{18}^{''})^{(2)}(T_{17},t)}$ are first augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{13}^{''})^{(1,1,)}(T_{14},t)}$, $\boxed{+(a_{14}^{''})^{(1,1,)}(T_{14},t)}$, $\boxed{+(a_{15}^{''})^{(1,1,)}(T_{14},t)}$ are second augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{20}^{''})^{(3,3,3)}(T_{21},t)}$, $\boxed{+(a_{21}^{''})^{(3,3,3)}(T_{21},t)}$, $\boxed{+(a_{22}^{''})^{(3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{24}^{''})^{(4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{25}^{''})^{(4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{26}^{''})^{(4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{28}^{''})^{(5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{29}^{''})^{(5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{30}^{''})^{(5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{32}^{''})^{(6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{33}^{''})^{(6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{34}^{''})^{(6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{36}^{''})^{(7,7,7)}(T_{37},t)}$, $\boxed{+(a_{37}^{''})^{(7,7,7)}(T_{37},t)}$, $\boxed{+(a_{38}^{''})^{(7,7,7)}(T_{37},t)}$ are seventh augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{40}^{''})^{(8,8,8)}(T_{41},t)}$, $\boxed{+(a_{41}^{''})^{(8,8,8)}(T_{41},t)}$, $\boxed{+(a_{42}^{''})^{(8,8,8)}(T_{41},t)}$ are eight augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{44}^{''})^{(9,9)}(T_{45},t)}$, $\boxed{+(a_{45}^{''})^{(9,9)}(T_{45},t)}$, $\boxed{+(a_{46}^{''})^{(9,9)}(T_{45},t)}$ are ninth augmentation coefficient for category 1, 2 and 3 | |
| $\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \left[\begin{array}{ccc} (b_{16}^{'})^{(2)}\boxed{-(b_{16}^{''})^{(2)}(G_{19},t)} & \boxed{-(b_{13}^{''})^{(1,1,)}(G,t)} & \boxed{-(b_{20}^{''})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b_{24}^{''})^{(4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{28}^{''})^{(5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{32}^{''})^{(6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{36}^{''})^{(7,7,7)}(G_{39},t)} & \boxed{-(b_{40}^{''})^{(8,8,8)}(G_{43},t)} & \boxed{-(b_{44}^{''})^{(9,9)}(G_{47},t)} \end{array}\right] T_{16}$ | 64 |
| $\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \left[\begin{array}{ccc} (b_{17}^{'})^{(2)}\boxed{-(b_{17}^{''})^{(2)}(G_{19},t)} & \boxed{-(b_{14}^{''})^{(1,1,)}(G,t)} & \boxed{-(b_{21}^{''})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b_{25}^{''})^{(4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{29}^{''})^{(5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{33}^{''})^{(6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{37}^{''})^{(7,7,7)}(G_{39},t)} & \boxed{-(b_{41}^{''})^{(8,8,8)}(G_{43},t)} & \boxed{-(b_{45}^{''})^{(9,9)}(G_{47},t)} \end{array}\right] T_{17}$ | 65 |
| $\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \left[\begin{array}{ccc} (b_{18}^{'})^{(2)}\boxed{-(b_{18}^{''})^{(2)}(G_{19},t)} & \boxed{-(b_{15}^{''})^{(1,1,)}(G,t)} & \boxed{-(b_{22}^{''})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b_{26}^{''})^{(4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{30}^{''})^{(5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{34}^{''})^{(6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{38}^{''})^{(7,7,7)}(G_{39},t)} & \boxed{-(b_{42}^{''})^{(8,8,8)}(G_{43},t)} & \boxed{-(b_{46}^{''})^{(9,9)}(G_{47},t)} \end{array}\right] T_{18}$ | 66 |
| where $\boxed{-(b_{16}^{''})^{(2)}(G_{19},t)}$, $\boxed{-(b_{17}^{''})^{(2)}(G_{19},t)}$, $\boxed{-(b_{18}^{''})^{(2)}(G_{19},t)}$ are first detrition coefficients for category 1, 2 and 3 | |

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| $-(b_{13}^{''})^{(1,1,)}(G,t)$, $-(b_{14}^{''})^{(1,1,)}(G,t)$, $-(b_{15}^{''})^{(1,1,)}(G,t)$ are second detrition coefficients for category 1,2 and 3<br>$-(b_{20}^{''})^{(3,3,3,)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,)}(G_{23},t)$ are third detrition coefficients for category 1,2 and 3<br>$-(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1,2 and 3<br>$-(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1,2 and 3<br>$-(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$ are sixth detrition coefficients for category 1,2 and 3<br>$-(b_{36}^{''})^{(7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1,2 and 3<br>$-(b_{40}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8)}(G_{43},t)$ are eight detrition coefficients for category 1,2 and 3<br>$-(b_{44}^{''})^{(9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1,2 and 3 | |
| $\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \left[ \begin{array}{c} (a_{20}^{'})^{(3)} +(a_{20}^{''})^{(3)}(T_{21},t) +(a_{16}^{''})^{(2,2,2)}(T_{17},t) +(a_{13}^{''})^{(1,1,1,)}(T_{14},t) \\ +(a_{24}^{''})^{(4,4,4,4,4,4)}(T_{25},t) +(a_{28}^{''})^{(5,5,5,5,5,5)}(T_{29},t) +(a_{32}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ +(a_{36}^{''})^{(7,7,7,7)}(T_{37},t) +(a_{40}^{''})^{(8,8,8,8)}(T_{41},t) +(a_{44}^{''})^{(9,9,9)}(T_{45},t) \end{array} \right] G_{20}$ | 67 |
| $\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \left[ \begin{array}{c} (a_{21}^{'})^{(3)} +(a_{21}^{''})^{(3)}(T_{21},t) +(a_{17}^{''})^{(2,2,2)}(T_{17},t) +(a_{14}^{''})^{(1,1,1,)}(T_{14},t) \\ +(a_{25}^{''})^{(4,4,4,4,4,4)}(T_{25},t) +(a_{29}^{''})^{(5,5,5,5,5,5)}(T_{29},t) +(a_{33}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ +(a_{37}^{''})^{(7,7,7,7)}(T_{37},t) +(a_{41}^{''})^{(8,8,8,8)}(T_{41},t) +(a_{45}^{''})^{(9,9,9)}(T_{45},t) \end{array} \right] G_{21}$ | 68 |
| $\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \left[ \begin{array}{c} (a_{22}^{'})^{(3)} +(a_{22}^{''})^{(3)}(T_{21},t) +(a_{18}^{''})^{(2,2,2)}(T_{17},t) +(a_{15}^{''})^{(1,1,1,)}(T_{14},t) \\ +(a_{26}^{''})^{(4,4,4,4,4,4)}(T_{25},t) +(a_{30}^{''})^{(5,5,5,5,5,5)}(T_{29},t) +(a_{34}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ +(a_{38}^{''})^{(7,7,7,7)}(T_{37},t) +(a_{42}^{''})^{(8,8,8,8)}(T_{41},t) +(a_{46}^{''})^{(9,9,9)}(T_{45},t) \end{array} \right] G_{22}$ | 69 |
| $+(a_{20}^{''})^{(3)}(T_{21},t)$, $+(a_{21}^{''})^{(3)}(T_{21},t)$, $+(a_{22}^{''})^{(3)}(T_{21},t)$ are first augmentation coefficients for category 1, 2 and 3<br>$+(a_{16}^{''})^{(2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2)}(T_{17},t)$ are second augmentation coefficients for category 1, 2 and 3<br>$+(a_{13}^{''})^{(1,1,1,)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,)}(T_{14},t)$ are third augmentation coefficients for category 1, 2 and 3<br>$+(a_{24}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficients for category 1, 2 and 3<br>$+(a_{28}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficients for category 1, 2 and 3<br>$+(a_{32}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation | |

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| coefficients for category 1, 2 and 3 $+(a_{36}^{''})^{(7,7,7,7)}(T_{37},t)$, $+(a_{37}^{''})^{(7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7)}(T_{37},t)$ are seventh augmentation coefficients for category 1, 2 and 3 $+(a_{40}^{''})^{(8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8)}(T_{41},t)$ are eight augmentation coefficients for category 1, 2 and 3 $+(a_{44}^{''})^{(9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9)}(T_{45},t)$, $+(a_{46}^{''})^{(9,9,9)}(T_{45},t)$ are ninth augmentation coefficients for category 1, 2 and 3 | |
| $\frac{dT_{20}}{dt} = (b_{20})^{(3)}T_{21} - \begin{bmatrix} (b_{20}^{'})^{(3)}\,-(b_{20}^{''})^{(3)}(G_{23},t) & -(b_{16}^{''})^{(2,2,2)}(G_{19},t) & -(b_{13}^{''})^{(1,1,1,)}(G,t) \\ -(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t) & -(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t) & -(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t) \\ -(b_{36}^{''})^{(7,7,7,7)}(G_{39},t) & -(b_{40}^{''})^{(8,8,8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9,9)}(G_{47},t) \end{bmatrix} T_{20}$ | *70* |
| $\frac{dT_{21}}{dt} = (b_{21})^{(3)}T_{20} - \begin{bmatrix} (b_{21}^{'})^{(3)}\,-(b_{21}^{''})^{(3)}(G_{23},t) & -(b_{17}^{''})^{(2,2,2)}(G_{19},t) & -(b_{14}^{''})^{(1,1,1,)}(G,t) \\ -(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t) & -(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t) & -(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t) \\ -(b_{37}^{''})^{(7,7,7,7)}(G_{39},t) & -(b_{41}^{''})^{(8,8,8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9,9)}(G_{47},t) \end{bmatrix} T_{21}$ | *71* |
| $\frac{dT_{22}}{dt} = (b_{22})^{(3)}T_{21} - \begin{bmatrix} (b_{22}^{'})^{(3)}\,-(b_{22}^{''})^{(3)}(G_{23},t) & -(b_{18}^{''})^{(2,2,2)}(G_{19},t) & -(b_{15}^{''})^{(1,1,1,)}(G,t) \\ -(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t) & -(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t) & -(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t) \\ -(b_{38}^{''})^{(7,7,7,7)}(G_{39},t) & -(b_{42}^{''})^{(8,8,8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9)}(G_{47},t) \end{bmatrix} T_{22}$ | *72* |
| $-(b_{20}^{''})^{(3)}(G_{23},t)$, $-(b_{21}^{''})^{(3)}(G_{23},t)$, $-(b_{22}^{''})^{(3)}(G_{23},t)$ *are first detrition coefficients for category 1, 2 and 3* $-(b_{16}^{''})^{(2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2)}(G_{19},t)$ *are second detrition coefficients for category 1, 2 and 3* $-(b_{13}^{''})^{(1,1,1,)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,)}(G,t)$, $-(b_{15}^{''})^{(1,1,1,)}(G,t)$ *are third detrition coefficients for category 1,2 and 3* $-(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$ *are fourth detrition coefficients for category 1, 2 and 3* $-(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$ *are fifth detrition coefficients for category 1, 2 and 3* $-(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$ *are sixth detrition coefficients for category 1, 2 and 3* $-(b_{36}^{''})^{(7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3 $-(b_{40}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8)}(G_{43},t)$ are eight detrition coefficients for category 1, 2 and 3 $-(b_{46}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3 | |

$$\frac{dG_{24}}{dt} = (a_{24})^{(4)} G_{25} - \begin{bmatrix} (a'_{24})^{(4)} & +(a''_{24})^{(4)}(T_{25},t) & +(a''_{28})^{(5,5,)}(T_{29},t) & +(a''_{32})^{(6,6,)}(T_{33},t) \\ & +(a''_{13})^{(1,1,1,1)}(T_{14},t) & +(a''_{16})^{(2,2,2,2)}(T_{17},t) & +(a''_{20})^{(3,3,3,3)}(T_{21},t) \\ & +(a''_{36})^{(7,7,7,7,7)}(T_{37},t) & +(a''_{40})^{(8,8,8,8,8)}(T_{41},t) & +(a''_{44})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{24} \quad (73)$$

$$\frac{dG_{25}}{dt} = (a_{25})^{(4)} G_{24} - \begin{bmatrix} (a'_{25})^{(4)} & +(a''_{25})^{(4)}(T_{25},t) & +(a''_{29})^{(5,5,)}(T_{29},t) & +(a''_{33})^{(6,6)}(T_{33},t) \\ & +(a''_{14})^{(1,1,1,1)}(T_{14},t) & +(a''_{17})^{(2,2,2,2)}(T_{17},t) & +(a''_{21})^{(3,3,3,3)}(T_{21},t) \\ & +(a''_{37})^{(7,7,7,7,7)}(T_{37},t) & +(a''_{41})^{(8,8,8,8,8)}(T_{41},t) & +(a''_{45})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{25} \quad (74)$$

$$\frac{dG_{26}}{dt} = (a_{26})^{(4)} G_{25} - \begin{bmatrix} (a'_{26})^{(4)} & +(a''_{26})^{(4)}(T_{25},t) & +(a''_{30})^{(5,5,)}(T_{29},t) & +(a''_{34})^{(6,6,)}(T_{33},t) \\ & +(a''_{15})^{(1,1,1,1)}(T_{14},t) & +(a''_{18})^{(2,2,2,2)}(T_{17},t) & +(a''_{22})^{(3,3,3,3)}(T_{21},t) \\ & +(a''_{38})^{(7,7,7,7,7)}(T_{37},t) & +(a''_{42})^{(8,8,8,8,8)}(T_{41},t) & +(a''_{46})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{26} \quad (75)$$

$(a''_{24})^{(4)}(T_{25},t)$, $(a''_{25})^{(4)}(T_{25},t)$, $(a''_{26})^{(4)}(T_{25},t)$ *are first augmentation coefficients category* 1, 2 3

$+(a''_{28})^{(5,5,)}(T_{29},t)$, $+(a''_{29})^{(5,5,)}(T_{29},t)$, $+(a''_{30})^{(5,5,)}(T_{29},t)$ *are second augmentation coefficient for category* 1, 2 *and* 3

$+(a''_{32})^{(6,6,)}(T_{33},t)$, $+(a''_{33})^{(6,6,)}(T_{33},t)$, $+(a''_{34})^{(6,6,)}(T_{33},t)$ *are third augmentation coefficient for category* 1, 2 *and* 3

$+(a''_{13})^{(1,1,1,1)}(T_{14},t)$, $+(a''_{14})^{(1,1,1,1)}(T_{14},t)$, $+(a''_{15})^{(1,1,1,1)}(T_{14},t)$ *are fourth augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{16})^{(2,2,2,2)}(T_{17},t)$, $+(a''_{17})^{(2,2,2,2)}(T_{17},t)$, $+(a''_{18})^{(2,2,2,2)}(T_{17},t)$ *are fifth augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{20})^{(3,3,3,3)}(T_{21},t)$, $+(a''_{21})^{(3,3,3,3)}(T_{21},t)$, $+(a''_{22})^{(3,3,3,3)}(T_{21},t)$ *are sixth augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{36})^{(7,7,7,7,7)}(T_{37},t)$, $+(a''_{37})^{(7,7,7,7,7)}(T_{37},t)$, $+(a''_{38})^{(7,7,7,7,7)}(T_{37},t)$ *are seventh augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{40})^{(8,8,8,8,8)}(T_{41},t)$, $+(a''_{41})^{(8,8,8,8,8)}(T_{41},t)$, $+(a''_{42})^{(8,8,8,8,8)}(T_{41},t)$ *are eighth augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{46})^{(9,9,9,9)}(T_{45},t)$, $+(a''_{45})^{(9,9,9,9)}(T_{45},t)$, $+(a''_{44})^{(9,9,9,9)}(T_{45},t)$ are ninth detrition coefficients for category 1 2 3

$$\frac{dT_{24}}{dt} = (b_{24})^{(4)} T_{25} - \begin{bmatrix} (b'_{24})^{(4)} & -(b''_{24})^{(4)}(G_{27},t) & -(b''_{28})^{(5,5,)}(G_{31},t) & -(b''_{32})^{(6,6,)}(G_{35},t) \\ & -(b''_{13})^{(1,1,1,1)}(G,t) & -(b''_{16})^{(2,2,2,2)}(G_{19},t) & -(b''_{20})^{(3,3,3,3)}(G_{23},t) \\ & -(b''_{36})^{(7,7,7,7,7)}(G_{39},t) & -(b''_{40})^{(8,8,8,8,8)}(G_{43},t) & -(b''_{44})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{24} \quad (76)$$

$$\frac{dT_{25}}{dt} = (b_{25})^{(4)} T_{24} - \begin{bmatrix} (b'_{25})^{(4)} & -(b''_{25})^{(4)}(G_{27},t) & -(b''_{29})^{(5,5,)}(G_{31},t) & -(b''_{33})^{(6,6,)}(G_{35},t) \\ & -(b''_{14})^{(1,1,1,1)}(G,t) & -(b''_{17})^{(2,2,2,2)}(G_{19},t) & -(b''_{21})^{(3,3,3,3)}(G_{23},t) \\ & -(b''_{37})^{(7,7,7,7,7)}(G_{39},t) & -(b''_{41})^{(8,8,8,8,8)}(G_{43},t) & -(b''_{45})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{25} \quad (77)$$

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| $\frac{dT_{26}}{dt} = (b_{26})^{(4)}T_{25} - \begin{bmatrix} (b'_{26})^{(4)}\boxed{-(b''_{26})^{(4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,)}(G_{35},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1)}(G,t)} & \boxed{-(b''_{18})^{(2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{38})^{(7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{26}$ | 78 |
| $Where$ $\boxed{-(b''_{24})^{(4)}(G_{27},t)}, \boxed{-(b''_{25})^{(4)}(G_{27},t)}, \boxed{-(b''_{26})^{(4)}(G_{27},t)}$ $are\ first\ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b''_{28})^{(5,5,)}(G_{31},t)}, \boxed{-(b''_{29})^{(5,5,)}(G_{31},t)}, \boxed{-(b''_{30})^{(5,5,)}(G_{31},t)}$ $are\ second\ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b''_{32})^{(6,6,)}(G_{35},t)}, \boxed{-(b''_{33})^{(6,6,)}(G_{35},t)}, \boxed{-(b''_{34})^{(6,6,)}(G_{35},t)}$ $are\ third\ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b''_{13})^{(1,1,1,1)}(G,t)}, \boxed{-(b''_{14})^{(1,1,1,1)}(G,t)}, \boxed{-(b''_{15})^{(1,1,1,1)}(G,t)}$ $are\ fourth\ \ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b''_{16})^{(2,2,2,2)}(G_{19},t)}, \boxed{-(b''_{17})^{(2,2,2,2)}(G_{19},t)}, \boxed{-(b''_{18})^{(2,2,2,2)}(G_{19},t)}$ $are\ fifth\ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b''_{20})^{(3,3,3,3)}(G_{23},t)}, \boxed{-(b''_{21})^{(3,3,3,3)}(G_{23},t)}, \boxed{-(b''_{22})^{(3,3,3,3)}(G_{23},t)}$ $are\ sixth\ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b''_{36})^{(7,7,7,7,7)}(G_{39},t)}, \boxed{-(b''_{37})^{(7,7,7,7,7)}(G_{39},t)}, \boxed{-(b''_{38})^{(7,7,7,7,7)}(G_{39},t)}$ $are\ seventh\ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b''_{40})^{(8,8,8,8,8)}(G_{43},t)}, \boxed{-(b''_{41})^{(8,8,8,8,8)}(G_{43},t)}, \boxed{-(b''_{42})^{(8,8,8,8,8)}(G_{43},t)}$ $are\ eighth\ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b''_{46})^{(9,9,9,9)}(G_{47},t)}, \boxed{-(b''_{45})^{(9,9,9,9)}(G_{47},t)}, \boxed{-(b''_{44})^{(9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1 2 3 | |
| $\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \begin{bmatrix} (a'_{28})^{(5)}\boxed{+(a''_{28})^{(5)}(T_{29},t)} & \boxed{+(a''_{24})^{(4,4,)}(T_{25},t)} & \boxed{+(a''_{32})^{(6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{36})^{(7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{40})^{(8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{44})^{(9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{28}$ | 79 |
| $\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \begin{bmatrix} (a'_{29})^{(5)}\boxed{+(a''_{29})^{(5)}(T_{29},t)} & \boxed{+(a''_{25})^{(4,4,)}(T_{25},t)} & \boxed{+(a''_{33})^{(6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{14})^{(1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{17})^{(2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{21})^{(3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{37})^{(7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{41})^{(8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{45})^{(9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{29}$ | 80 |
| $\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \begin{bmatrix} (a'_{30})^{(5)}\boxed{+(a''_{30})^{(5)}(T_{29},t)} & \boxed{+(a''_{26})^{(4,4,)}(T_{25},t)} & \boxed{+(a''_{34})^{(6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{15})^{(1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{18})^{(2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{22})^{(3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{38})^{(7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{42})^{(8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{46})^{(9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{30}$ | 81 |

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| *Where* $+(a_{28}^{''})^{(5)}(T_{29},t)$ , $+(a_{29}^{''})^{(5)}(T_{29},t)$ , $+(a_{30}^{''})^{(5)}(T_{29},t)$ *are first augmentation coefficients for category* 1, 2 *and* 3<br>*And* $+(a_{24}^{''})^{(4,4,)}(T_{25},t)$ , $+(a_{25}^{''})^{(4,4,)}(T_{25},t)$ , $+(a_{26}^{''})^{(4,4,)}(T_{25},t)$ *are second augmentation coefficient for category* 1, 2 *and* 3<br>$+(a_{32}^{''})^{(6,6,6)}(T_{33},t)$ , $+(a_{33}^{''})^{(6,6,6)}(T_{33},t)$ , $+(a_{34}^{''})^{(6,6,6)}(T_{33},t)$ *are third augmentation coefficient for category* 1, 2 *and* 3<br>$+(a_{13}^{''})^{(1,1,1,1,1)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1,1)}(T_{14},t)$ are fourth augmentation coefficients for category 1,2, and 3<br>$+(a_{16}^{''})^{(2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2)}(T_{17},t)$ are fifth augmentation coefficients for category 1,2,and 3<br>$+(a_{20}^{''})^{(3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3)}(T_{21},t)$ are sixth augmentation coefficients for category 1,2, 3<br>$+(a_{36}^{''})^{(7,7,7,7,7,7)}(T_{37},t)$ , $+(a_{37}^{''})^{(7,7,7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7,7,7)}(T_{37},t)$ are seventh augmentation coefficients for category 1,2, 3<br>$+(a_{40}^{''})^{(8,8\ ,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8,8,8)}(T_{41},t)$ are eighth augmentation coefficients for category 1,2, 3<br>$+(a_{46}^{''})^{(9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficients for category 1,2, 3 | |
| $\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \left[\begin{matrix}(b_{28}^{'})^{(5)}\boxed{-(b_{28}^{''})^{(5)}(G_{31},t)}\ \boxed{-(b_{24}^{''})^{(4,4,)}(G_{27},t)}\ \boxed{-(b_{32}^{''})^{(6,6,6)}(G_{35},t)} \\ \boxed{-(b_{13}^{''})^{(1,1,1,1,1)}(G,t)}\ \boxed{-(b_{16}^{''})^{(2,2,2,2,2)}(G_{19},t)}\ \boxed{-(b_{20}^{''})^{(3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{36}^{''})^{(7,7,7,7,7,7)}(G_{39},t)}\ \boxed{-(b_{40}^{''})^{(8,8,8,8,8,8)}(G_{43},t)}\ \boxed{-(b_{44}^{''})^{(9,9,9,9,9)}(G_{47},t)}\end{matrix}\right] T_{28}$ | *82* |
| $\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \left[\begin{matrix}(b_{29}^{'})^{(5)}\boxed{-(b_{29}^{''})^{(5)}(G_{31},t)}\ \boxed{-(b_{25}^{''})^{(4,4,)}(G_{27},t)}\ \boxed{-(b_{33}^{''})^{(6,6,6)}(G_{35},t)} \\ \boxed{-(b_{14}^{''})^{(1,1,1,1,1)}(G,t)}\ \boxed{-(b_{17}^{''})^{(2,2,2,2,2)}(G_{19},t)}\ \boxed{-(b_{21}^{''})^{(3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{37}^{''})^{(7,7,7,7,7,7)}(G_{39},t)}\ \boxed{-(b_{41}^{''})^{(8,8,8,8,8,8)}(G_{43},t)}\ \boxed{-(b_{45}^{''})^{(9,9,9,9,9)}(G_{47},t)}\end{matrix}\right] T_{29}$ | *83* |
| $\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \left[\begin{matrix}(b_{30}^{'})^{(5)}\boxed{-(b_{30}^{''})^{(5)}(G_{31},t)}\ \boxed{-(b_{26}^{''})^{(4,4,)}(G_{27},t)}\ \boxed{-(b_{34}^{''})^{(6,6,6)}(G_{35},t)} \\ \boxed{-(b_{15}^{''})^{(1,1,1,1,1,)}(G,t)}\ \boxed{-(b_{18}^{''})^{(2,2,2,2,2)}(G_{19},t)}\ \boxed{-(b_{22}^{''})^{(3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{38}^{''})^{(7,7,7,7,7,7)}(G_{39},t)}\ \boxed{-(b_{42}^{''})^{(8,8,8,8,8,8)}(G_{43},t)}\ \boxed{-(b_{46}^{''})^{(9,9,9,9,9)}(G_{47},t)}\end{matrix}\right] T_{30}$ | *84* |
| *where* $\boxed{-(b_{28}^{''})^{(5)}(G_{31},t)}$ , $\boxed{-(b_{29}^{''})^{(5)}(G_{31},t)}$ , $\boxed{-(b_{30}^{''})^{(5)}(G_{31},t)}$ *are first detrition coefficients for category* 1, 2 *and* 3<br>$\boxed{-(b_{24}^{''})^{(4,4,)}(G_{27},t)}$, $\boxed{-(b_{25}^{''})^{(4,4,)}(G_{27},t)}$, $\boxed{-(b_{26}^{''})^{(4,4,)}(G_{27},t)}$ *are second detrition coefficients for category* 1,2 *and* 3<br>$\boxed{-(b_{32}^{''})^{(6,6,6)}(G_{35},t)}$, $\boxed{-(b_{33}^{''})^{(6,6,6)}(G_{35},t)}$, $\boxed{-(b_{34}^{''})^{(6,6,6)}(G_{35},t)}$ *are third detrition coefficients for category* 1,2 *and* 3<br>$\boxed{-(b_{13}^{''})^{(1,1,1,1,1)}(G,t)}$, $\boxed{-(b_{14}^{''})^{(1,1,1,1,1)}(G,t)}$ , $\boxed{-(b_{15}^{''})^{(1,1,1,1,1,)}(G,t)}$ are fourth detrition coefficients for category 1,2, and 3<br>$\boxed{-(b_{16}^{''})^{(2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b_{17}^{''})^{(2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b_{18}^{''})^{(2,2,2,2,2)}(G_{19},t)}$ are fifth detrition coefficients | |

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| for category 1,2, and 3<br>$\boxed{-(b_{20}^{''})^{(3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b_{21}^{''})^{(3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b_{22}^{''})^{(3,3,3,3,3)}(G_{23},t)}$ are sixth detrition coefficients for category 1,2, and 3<br>$\boxed{-(b_{36}^{''})^{(7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{37}^{''})^{(7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{38}^{''})^{(7,7,7,7,7,7)}(G_{39},t)}$ are seventh detrition coefficients for category 1,2, and 3<br>$\boxed{-(b_{42}^{''})^{(8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{41}^{''})^{(8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{40}^{''})^{(8,8,8,8,8,8)}(G_{43},t)}$ are eighth detrition coefficients for category 1,2, and 3<br>$\boxed{-(b_{46}^{''})^{(9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b_{45}^{''})^{(9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b_{44}^{''})^{(9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1,2, and 3 | |
| $\frac{dG_{32}}{dt} = (a_{32})^{(6)}G_{33} - \left[\begin{matrix} (a_{32}^{'})^{(6)}\boxed{+(a_{32}^{''})^{(6)}(T_{33},t)} & \boxed{+(a_{28}^{''})^{(5,5,5)}(T_{29},t)} & \boxed{+(a_{24}^{''})^{(4,4,4,)}(T_{25},t)} \\ \boxed{+(a_{13}^{''})^{(1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{16}^{''})^{(2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{20}^{''})^{(3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{36}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{40}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{44}^{''})^{(9,9,9,9,9,9)}(T_{45},t)} \end{matrix}\right] G_{32}$ | *85* |
| $\frac{dG_{33}}{dt} = (a_{33})^{(6)}G_{32} - \left[\begin{matrix} (a_{33}^{'})^{(6)}\boxed{+(a_{33}^{''})^{(6)}(T_{33},t)} & \boxed{+(a_{29}^{''})^{(5,5,5)}(T_{29},t)} & \boxed{+(a_{25}^{''})^{(4,4,4,)}(T_{25},t)} \\ \boxed{+(a_{14}^{''})^{(1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{17}^{''})^{(2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{21}^{''})^{(3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{37}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{41}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{45}^{''})^{(9,9,9,9,9,9)}(T_{45},t)} \end{matrix}\right] G_{33}$ | *86* |
| $\frac{dG_{34}}{dt} = (a_{34})^{(6)}G_{33} - \left[\begin{matrix} (a_{34}^{'})^{(6)}\boxed{+(a_{34}^{''})^{(6)}(T_{33},t)} & \boxed{+(a_{30}^{''})^{(5,5,5)}(T_{29},t)} & \boxed{+(a_{26}^{''})^{(4,4,4,)}(T_{25},t)} \\ \boxed{+(a_{15}^{''})^{(1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{18}^{''})^{(2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{22}^{''})^{(3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{38}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{42}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{46}^{''})^{(9,9,9,9,9,9)}(T_{45},t)} \end{matrix}\right] G_{34}$ | *87* |
| $\boxed{+(a_{32}^{''})^{(6)}(T_{33},t)}$, $\boxed{+(a_{33}^{''})^{(6)}(T_{33},t)}$, $\boxed{+(a_{34}^{''})^{(6)}(T_{33},t)}$ *are first augmentation coefficients for category* 1, 2 *and* 3<br>$\boxed{+(a_{28}^{''})^{(5,5,5)}(T_{29},t)}$, $\boxed{+(a_{29}^{''})^{(5,5,5)}(T_{29},t)}$, $\boxed{+(a_{30}^{''})^{(5,5,5)}(T_{29},t)}$ *are second augmentation coefficients for category* 1, 2 *and* 3<br>$\boxed{+(a_{24}^{''})^{(4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{25}^{''})^{(4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{26}^{''})^{(4,4,4,)}(T_{25},t)}$ *are third augmentation coefficients for category* 1, 2 *and* 3<br>$\boxed{+(a_{13}^{''})^{(1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a_{14}^{''})^{(1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a_{15}^{''})^{(1,1,1,1,1,1)}(T_{14},t)}$ - are fourth augmentation coefficients<br>$\boxed{+(a_{16}^{''})^{(2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a_{17}^{''})^{(2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a_{18}^{''})^{(2,2,2,2,2,2)}(T_{17},t)}$ - fifth augmentation coefficients<br>$\boxed{+(a_{20}^{''})^{(3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a_{21}^{''})^{(3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a_{22}^{''})^{(3,3,3,3,3,3)}(T_{21},t)}$ sixth augmentation coefficients<br>$\boxed{+(a_{36}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)}$, $\boxed{+(a_{37}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)}$, $\boxed{+(a_{38}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)}$ seventh augmentation coefficients<br>$\boxed{+(a_{40}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{41}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{42}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)}$ Eighth augmentation coefficients<br>$\boxed{+(a_{44}^{''})^{(9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{45}^{''})^{(9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{46}^{''})^{(9,9,9,9,9,9)}(T_{45},t)}$ ninth augmentation | |

coefficients

$$\frac{dT_{32}}{dt} = (b_{32})^{(6)}T_{33} - \begin{bmatrix} (b'_{32})^{(6)}\boxed{-(b''_{32})^{(6)}(G_{35},t)} & \boxed{-(b''_{28})^{(5,5,5)}(G_{31},t)} & \boxed{-(b''_{24})^{(4,4,4,)}(G_{27},t)} \\ \boxed{-(b''_{13})^{(1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{16})^{(2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{20})^{(3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{36})^{(7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{40})^{(8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{44})^{(9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{32} \quad (88)$$

$$\frac{dT_{33}}{dt} = (b_{33})^{(6)}T_{32} - \begin{bmatrix} (b'_{33})^{(6)}\boxed{-(b''_{33})^{(6)}(G_{35},t)} & \boxed{-(b''_{29})^{(5,5,5)}(G_{31},t)} & \boxed{-(b''_{25})^{(4,4,4,)}(G_{27},t)} \\ \boxed{-(b''_{14})^{(1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{17})^{(2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{21})^{(3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{37})^{(7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{41})^{(8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{45})^{(9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{33} \quad (89)$$

$$\frac{dT_{34}}{dt} = (b_{34})^{(6)}T_{33} - \begin{bmatrix} (b'_{34})^{(6)}\boxed{-(b''_{34})^{(6)}(G_{35},t)} & \boxed{-(b''_{30})^{(5,5,5)}(G_{31},t)} & \boxed{-(b''_{26})^{(4,4,4,)}(G_{27},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{18})^{(2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{38})^{(7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{34} \quad (90)$$

$\boxed{-(b''_{32})^{(6)}(G_{35},t)}, \boxed{-(b''_{33})^{(6)}(G_{35},t)}, \boxed{-(b''_{34})^{(6)}(G_{35},t)}$ *are first detrition coefficients for category* 1, 2 *and* 3

$\boxed{-(b''_{28})^{(5,5,5)}(G_{31},t)}, \boxed{-(b''_{29})^{(5,5,5)}(G_{31},t)}, \boxed{-(b''_{30})^{(5,5,5)}(G_{31},t)}$ *are second detrition coefficients for category* 1, 2 *and* 3

$\boxed{-(b''_{24})^{(4,4,4,)}(G_{27},t)}, \boxed{-(b''_{25})^{(4,4,4,)}(G_{27},t)}, \boxed{-(b''_{26})^{(4,4,4,)}(G_{27},t)}$ *are third detrition coefficients for category* 1,2 *and* 3

$\boxed{-(b''_{13})^{(1,1,1,1,1,1)}(G,t)}, \boxed{-(b''_{14})^{(1,1,1,1,1,1)}(G,t)}, \boxed{-(b''_{15})^{(1,1,1,1,1,1)}(G,t)}$ are fourth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{16})^{(2,2,2,2,2,2)}(G_{19},t)}, \boxed{-(b''_{17})^{(2,2,2,2,2,2)}(G_{19},t)}, \boxed{-(b''_{18})^{(2,2,2,2,2,2)}(G_{19},t)}$ are fifth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{20})^{(3,3,3,3,3,3)}(G_{23},t)}, \boxed{-(b''_{21})^{(3,3,3,3,3,3)}(G_{23},t)}, \boxed{-(b''_{22})^{(3,3,3,3,3,3)}(G_{23},t)}$ are sixth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{36})^{(7,7,7,7,7,7,7)}(G_{39},t)}, \boxed{-(b''_{37})^{(7,7,7,7,7,7,7)}(G_{39},t)}, \boxed{-(b''_{38})^{(7,7,7,7,7,7,7)}(G_{39},t)}$ are seventh detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{40})^{(8,8,8,8,8,8,8)}(G_{43},t)}, \boxed{-(b''_{41})^{(8,8,8,8,8,8,8)}(G_{43},t)}, \boxed{-(b''_{42})^{(8,8,8,8,8,8,8)}(G_{43},t)}$ are eighth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{46})^{(9,9,9,9,9,9)}(G_{47},t)}, \boxed{-(b''_{45})^{(9,9,9,9,9,9)}(G_{47},t)}, \boxed{-(b''_{44})^{(9,9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2, and 3

$$\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \begin{bmatrix} (a'_{36})^{(7)}\boxed{+(a''_{36})^{(7)}(T_{37},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{24})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8,)}(T_{41},t)} & \boxed{+(a''_{44})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{13} \quad (91)$$

$$\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \begin{bmatrix} (a_{37}')^{(7)}\boxed{+(a_{37}'')^{(7)}(T_{37},t)} & \boxed{+(a_{17}'')^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{21}'')^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{25}'')^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{29}'')^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{33}'')^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{13}'')^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{41}'')^{(8,8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{45}'')^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{14} \quad (92)$$

$$\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \begin{bmatrix} (a_{38}')^{(7)}\boxed{+(a_{38}'')^{(7)}(T_{37},t)} & \boxed{+(a_{18}'')^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{22}'')^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{26}'')^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{30}'')^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{34}'')^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{15}'')^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{42}'')^{(8,8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{46}'')^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{15} \quad (93)$$

Where $\boxed{(a_{36}'')^{(7)}(T_{37},t)}$, $\boxed{(a_{37}'')^{(7)}(T_{37},t)}$, $\boxed{(a_{38}'')^{(7)}(T_{37},t)}$ are first augmentation coefficients for category 1, 2 and 3

$\boxed{+(a_{16}'')^{(2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a_{17}'')^{(2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a_{18}'')^{(2,2,2,2,2,2,2)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{20}'')^{(3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a_{21}'')^{(3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a_{22}'')^{(3,3,3,3,3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{24}'')^{(4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{25}'')^{(4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{26}'')^{(4,4,4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{28}'')^{(5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{29}'')^{(5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{30}'')^{(5,5,5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{32}'')^{(6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{33}'')^{(6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{34}'')^{(6,6,6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{13}'')^{(1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a_{13}'')^{(1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a_{15}'')^{(1,1,1,1,1,1,1)}(T_{14},t)}$ are seventh augmentation coefficient for category 1, 2 and 3

$\boxed{+(a_{42}'')^{(8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{41}'')^{(8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{40}'')^{(8,8,8,8,8,8,8,8,)}(T_{41},t)}$ are eighth augmentation coefficient for 1,2,3

$\boxed{+(a_{46}'')^{(9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{45}'')^{(9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{44}'')^{(9,9,9,9,9,9,9)}(T_{45},t)}$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{36}}{dt} = (b_{36})^{(7)}T_{37} - \begin{bmatrix} (b_{36}')^{(7)}\boxed{-(b_{36}'')^{(7)}(G_{39},t)} & \boxed{-(b_{16}'')^{(2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{20}'')^{(3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{24}'')^{(4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{28}'')^{(5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{32}'')^{(6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{13}'')^{(1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{40}'')^{(8,8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b_{44}'')^{(9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{13} \quad (94)$$

$$\frac{dT_{37}}{dt} =$$

$$(b_{37})^{(7)}T_{36} - \begin{bmatrix} (b'_{37})^{(7)}\boxed{-(b''_{37})^{(7)}(G_{39},t)} & \boxed{-(b''_{17})^{(2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{21})^{(3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{25})^{(4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{29})^{(5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{33})^{(6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{14})^{(1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{41})^{(8,8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{45})^{(9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{14}$$

$$\frac{dT_{38}}{dt} =$$

$$(b_{38})^{(7)}T_{37} - \begin{bmatrix} (b'_{38})^{(7)}\boxed{-(b''_{38})^{(7)}(G_{39},t)} & \boxed{-(b''_{18})^{(2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{26})^{(4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{42})^{(8,8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{15}$$

Where $\boxed{-(b''_{36})^{(7)}(G_{39},t)}$, $\boxed{-(b''_{37})^{(7)}(G_{39},t)}$, $\boxed{-(b''_{38})^{(7)}(G_{39},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{16})^{(2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{17})^{(2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{18})^{(2,2,2,2,2,2,2)}(G_{19},t)}$ are second detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{20})^{(3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{21})^{(3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{22})^{(3,3,3,3,3,3,3)}(G_{23},t)}$ are third detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{24})^{(4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{25})^{(4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{26})^{(4,4,4,4,4,4,4)}(G_{27},t)}$ are fourth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{28})^{(5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{29})^{(5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{30})^{(5,5,5,5,5,5,5)}(G_{31},t)}$ are fifth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{32})^{(6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b''_{33})^{(6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b''_{34})^{(6,6,6,6,6,6,6)}(G_{35},t)}$ are sixth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{15})^{(1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{14})^{(1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{13})^{(1,1,1,1,1,1,1)}(G,t)}$ are seventh detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{40})^{(8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{41})^{(8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{42})^{(8,8,8,8,8,8,8,8)}(G_{43},t)}$ are eighth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{46})^{(9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{45})^{(9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{44})^{(9,9,9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2 and 3

$$\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \begin{bmatrix} (a'_{40})^{(8)}\boxed{+(a''_{40})^{(8)}(T_{41},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{24})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{36})^{(7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{44})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{13}$$

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$$\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \begin{bmatrix} (a'_{41})^{(8)}\boxed{+(a''_{41})^{(8)}(T_{41},t)} & \boxed{+(a''_{17})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{21})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{25})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{29})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{33})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{37})^{(7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{45})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{14}$$

$$\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \begin{bmatrix} (a'_{42})^{(8)}\boxed{+(a''_{42})^{(8)}(T_{41},t)} & \boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{46})^{(9,9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{15}$$

Where $\boxed{+(a''_{40})^{(8)}(T_{41},t)}$, $\boxed{+(a''_{41})^{(8)}(T_{41},t)}$, $\boxed{+(a''_{42})^{(8)}(T_{41},t)}$ are first augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{16})^{(2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{17})^{(2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{20})^{(3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{21})^{(3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{24})^{(4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{25})^{(4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{28})^{(5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{29})^{(5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{32})^{(6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{33})^{(6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{13})^{(1,1,1,1,1,1,1,1)}(T_{14},t)}$ $\boxed{+(a''_{14})^{(1,1,1,1,1,1,1,1)}(T_{14},t)}$ $\boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1)}(T_{14},t)}$ are seventh augmentation coefficient for 1,2,3

$\boxed{+(a''_{36})^{(7,7,7,7,7,7,7,7)}(T_{37},t)}$, $\boxed{+(a''_{37})^{(7,7,7,7,7,7,7,7)}(T_{37},t)}$, $\boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7)}(T_{37},t)}$ are eighth augmentation coefficient for 1,2,3

$\boxed{+(a''_{46})^{(9,9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a''_{45})^{(9,9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a''_{44})^{(9,9,9,9,9,9,9,9)}(T_{45},t)}$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{40}}{dt} = (b_{40})^{(8)}T_{41} - \begin{bmatrix} (b'_{40})^{(8)}\boxed{-(b''_{40})^{(8)}(G_{43},t)} & \boxed{-(b''_{16})^{(2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{20})^{(3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{24})^{(4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{28})^{(5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{32})^{(6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{13})^{(1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{36})^{(7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{44})^{(9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{13}$$

$$\frac{dT_{41}}{dt} = (b_{41})^{(8)}T_{40} - \begin{bmatrix} (b'_{41})^{(8)}\boxed{-(b''_{41})^{(8)}(G_{43},t)} & \boxed{-(b''_{17})^{(2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{21})^{(3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{25})^{(4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{29})^{(5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{33})^{(6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{14})^{(1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{37})^{(7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{45})^{(9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{14}$$

$$\frac{dT_{42}}{dt} =$$

$$(b_{42})^{(8)}T_{41} - \begin{bmatrix} (b'_{42})^{(8)}\boxed{-(b''_{42})^{(8)}(G_{43},t)} & \boxed{-(b''_{18})^{(2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{26})^{(4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{38})^{(7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{46})^{(9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{15}$$

Where $\boxed{-(b''_{36})^{(7)}(G_{39},t)}$, $\boxed{-(b''_{37})^{(7)}(G_{39},t)}$, $\boxed{-(b''_{38})^{(7)}(G_{39},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{16})^{(2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{17})^{(2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{18})^{(2,2,2,2,2,2,2,2)}(G_{19},t)}$ are second detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{20})^{(3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{21})^{(3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{22})^{(3,3,3,3,3,3,3,3)}(G_{23},t)}$ are third detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{24})^{(4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{25})^{(4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{26})^{(4,4,4,4,4,4,4,4)}(G_{27},t)}$ are fourth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{28})^{(5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{29})^{(5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{30})^{(5,5,5,5,5,5,5,5)}(G_{31},t)}$ are fifth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{32})^{(6,6,6,6,)}(G_{35},t)}$, $\boxed{-(b''_{33})^{(6,6,6,6,)}(G_{35},t)}$, $\boxed{-(b''_{15})^{(1,1,1,1,1,1,1,1)}(G,t)}$ are sixth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{13})^{(1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{14})^{(1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{38})^{(7,7,)}(G_{39},t)}$ are seventh detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{36})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{37})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{38})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$ are eighth detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{44})^{(9,9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{45})^{(9,9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{46})^{(9,9,9,9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2 and 3

$$\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \begin{bmatrix} (a'_{44})^{(9)}\boxed{+(a''_{44})^{(9)}(T_{45},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{24})^{(4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{36})^{(7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8)}(T_{41},t)} \end{bmatrix} G_{13}$$

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$$\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \begin{bmatrix} (a'_{45})^{(9)}\boxed{+(a''_{45})^{(9)}(T_{45},t)} & \boxed{+(a''_{17})^{(2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{21})^{(3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{25})^{(4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{29})^{(5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{33})^{(6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{14})^{(1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{37})^{(7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{41})^{(8,8,8,8,8,8,8,8)}(T_{41},t)} \end{bmatrix} G_{14}$$

$$\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \begin{bmatrix} (a'_{46})^{(9)}\boxed{+(a''_{46})^{(9)}(T_{37},t)} & \boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{42})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)} \end{bmatrix} G_{15}$$

Where $\boxed{+(a''_{44})^{(9)}(T_{45},t)}$, $\boxed{+(a''_{45})^{(9)}(T_{45},t)}$, $\boxed{+(a''_{46})^{(9)}(T_{37},t)}$ are first augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{16})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{17})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{20})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{21})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{24})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{25})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{28})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{29})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{32})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{33})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{13})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a''_{14})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)}$ are Seventh augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)}$ $\boxed{+(a''_{37})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)}$ $\boxed{+(a''_{36})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)}$ are eighth augmentation coefficient for 1,2,3

$\boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a''_{42})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a''_{41})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)}$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{44}}{dt} = (b_{44})^{(9)}T_{45} - \begin{bmatrix} (b'_{44})^{(9)}\boxed{-(b''_{44})^{(9)}(G_{47},t)} & \boxed{-(b''_{16})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{20})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{24})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{28})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{32})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{13})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{36})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{40})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{bmatrix} T_{13}$$

$$\frac{dT_{45}}{dt} = (b_{45})^{(9)}T_{44} - \begin{bmatrix} (b'_{45})^{(9)}\boxed{-(b''_{45})^{(9)}(G_{47},t)} & \boxed{-(b''_{17})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{21})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{25})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{29})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{33})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{14})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{37})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{41})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{bmatrix} T_{14}$$

$$\frac{dT_{46}}{dt} =$$

$$(b_{46})^{(9)}T_{45} - \begin{bmatrix} (b_{46}^{'})^{(9)}\boxed{-(b_{46}^{''})^{(9)}(G_{47},t)} & \boxed{-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{bmatrix} T_{15}$$

Where $\boxed{-(b_{44}^{''})^{(9)}(G_{47},t)}$, $\boxed{-(b_{45}^{''})^{(9)}(G_{47},t)}$, $\boxed{-(b_{46}^{''})^{(9)}(G_{47},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)}$ are second detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)}$ are third detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)}$ are fourth detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)}$ are fifth detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)}$ are sixth detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)}$ are seventh detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)}$ are eighth detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$ are ninth detrition coefficients for category 1, 2 and 3

Where we suppose

$(a_i)^{(1)}, (a_i^{'})^{(1)}, (a_i^{''})^{(1)}, (b_i)^{(1)}, (b_i^{'})^{(1)}, (b_i^{''})^{(1)} > 0, \quad i,j = 13,14,15$ *97*

The functions $(a_i^{''})^{(1)}, (b_i^{''})^{(1)}$ are positive continuous increasing and bounded.

**Definition of** $(p_i)^{(1)}$, $(r_i)^{(1)}$:

$(a_i^{''})^{(1)}(T_{14},t) \le (p_i)^{(1)} \le (\hat{A}_{13})^{(1)}$

$(b_i^{''})^{(1)}(G,t) \le (r_i)^{(1)} \le (b_i^{'})^{(1)} \le (\hat{B}_{13})^{(1)}$

$lim_{T_2\to\infty}(a_i^{''})^{(1)}(T_{14},t) = (p_i)^{(1)}$ 98

$\lim_{G\to\infty}(b_i^{''})^{(1)}(G,t) = (r_i)^{(1)}$

**Definition of** $(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}$ :

Where $\boxed{(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}}$ are positive constants and $\boxed{i = 13,14,15}$

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|---|---|
| They satisfy Lipschitz condition:<br>$\|(a_i'')^{(1)}(T_{14}', t) - (a_i'')^{(1)}(T_{14}, t)\| \leq (\hat{k}_{13})^{(1)}\|T_{14} - T_{14}'\|e^{-(\hat{M}_{13})^{(1)}t}$<br>$\|(b_i'')^{(1)}(G', t) - (b_i'')^{(1)}(G, t)\| < (\hat{k}_{13})^{(1)}\|G - G'\|e^{-(\hat{M}_{13})^{(1)}t}$ | 99 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(1)}(T_{14}', t)$ and $(a_i'')^{(1)}(T_{14}, t)$ . $(T_{14}', t)$ and $(T_{14}, t)$ are points belonging to the interval $[(\hat{k}_{13})^{(1)}, (\hat{M}_{13})^{(1)}]$. It is to be noted that $(a_i'')^{(1)}(T_{14}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{13})^{(1)} = 1$ then the function $(a_i'')^{(1)}(T_{14}, t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}$ :<br>$(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)},$ are positive constants<br>$\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}} < 1$ | 100 |
| **Definition of** $(\hat{P}_{13})^{(1)}, (\hat{Q}_{13})^{(1)}$ :<br>There exists two constants $(\hat{P}_{13})^{(1)}$ and $(\hat{Q}_{13})^{(1)}$ which together With $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}, (\hat{A}_{13})^{(1)}$ and $(\hat{B}_{13})^{(1)}$ and the constants $(a_i)^{(1)}, (a_i')^{(1)}, (b_i)^{(1)}, (b_i')^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}, i = 13,14,15,$ satisfy the inequalities<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(a_i)^{(1)} + (a_i')^{(1)} + (\hat{A}_{13})^{(1)} + (\hat{P}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(b_i)^{(1)} + (b_i')^{(1)} + (\hat{B}_{13})^{(1)} + (\hat{Q}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$ | 101 |
| Where we suppose | |
| $(a_i)^{(2)}, (a_i')^{(2)}, (a_i'')^{(2)}, (b_i)^{(2)}, (b_i')^{(2)}, (b_i'')^{(2)} > 0, \quad i, j = 16,17,18$ | |
| The functions $(a_i'')^{(2)}, (b_i'')^{(2)}$ are positive continuous increasing and bounded. | |
| **Definition of** $(p_i)^{(2)}$, $(r_i)^{(2)}$: | |
| $(a_i'')^{(2)}(T_{17}, t) \leq (p_i)^{(2)} \leq (\hat{A}_{16})^{(2)}$ | 102 |
| $(b_i'')^{(2)}(G_{19}, t) \leq (r_i)^{(2)} \leq (b_i')^{(2)} \leq (\hat{B}_{16})^{(2)}$ | 103 |
| $\lim_{T_2 \to \infty}(a_i'')^{(2)}(T_{17}, t) = (p_i)^{(2)}$ | 104 |
| $\lim_{G \to \infty}(b_i'')^{(2)}((G_{19}), t) = (r_i)^{(2)}$ | 105 |
| **Definition of** $(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}$ : | 106 |

| | |
|---|---|
| Where $\boxed{(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}, (p_i)^{(2)}, \ (r_i)^{(2)}}$ are positive constants and $\boxed{i = 16{,}17{,}18}$ | |
| They satisfy Lipschitz condition: | |
| $\lvert(a_i^{''})^{(2)}(T_{17}^{'}, t) - (a_i^{''})^{(2)}(T_{17}, t)\rvert \leq (\hat{k}_{16})^{(2)} \lvert T_{17} - T_{17}^{'}\rvert e^{-(\hat{M}_{16})^{(2)}t}$ | 107 |
| $\lvert(b_i^{''})^{(2)}((G_{19})^{'}, t) - (b_i^{''})^{(2)}\big((G_{19}), t\big)\rvert < (\hat{k}_{16})^{(2)} \lvert\lvert(G_{19}) - (G_{19})^{'}\rvert\rvert e^{-(\hat{M}_{16})^{(2)}t}$ | 108 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(2)}(T_{17}^{'}, t)$ and $(a_i^{''})^{(2)}(T_{17}, t)$ . $(T_{17}^{'}, t)$ and $(T_{17}, t)$ are points belonging to the interval $\big[(\hat{k}_{16})^{(2)}, (\hat{M}_{16})^{(2)}\big]$. It is to be noted that $(a_i^{''})^{(2)}(T_{17}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{16})^{(2)} = 1$ then the function $(a_i^{''})^{(2)}(T_{17}, t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$ : | |
| $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$, are positive constants<br><br>$\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}$ , $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ | 109 |
| **Definition of** $(\hat{P}_{13})^{(2)}, (\hat{Q}_{13})^{(2)}$ :<br><br>There exists two constants $(\hat{P}_{16})^{(2)}$ and $(\hat{Q}_{16})^{(2)}$ which together with $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}, (\hat{A}_{16})^{(2)} and\ (\hat{B}_{16})^{(2)}$ and the constants $(a_i)^{(2)}, (a_i^{'})^{(2)}, (b_i)^{(2)}, (b_i^{'})^{(2)}, (p_i)^{(2)}, \ (r_i)^{(2)}, i = 16{,}17{,}18,$<br><br>satisfy the inequalities | |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(a_i)^{(2)} + (a_i^{'})^{(2)} + (\hat{A}_{16})^{(2)} + (\hat{P}_{16})^{(2)} (\hat{k}_{16})^{(2)}] < 1$ | 110 |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(b_i)^{(2)} + (b_i^{'})^{(2)} + (\hat{B}_{16})^{(2)} + (\hat{Q}_{16})^{(2)} (\hat{k}_{16})^{(2)}] < 1$ | 111 |
| Where we suppose | |
| $(a_i)^{(3)}, (a_i^{'})^{(3)}, (a_i^{''})^{(3)}, (b_i)^{(3)}, (b_i^{'})^{(3)}, (b_i^{''})^{(3)} > 0, \quad i, j = 20{,}21{,}22$<br><br>The functions $(a_i^{''})^{(3)}, (b_i^{''})^{(3)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(3)}, \ (\mathrm{r_i})^{(3)}$:<br><br>$(a_i^{''})^{(3)}(T_{21}, t) \leq (p_i)^{(3)} \leq (\hat{A}_{20})^{(3)}$<br><br>$(b_i^{''})^{(3)}(G_{23}, t) \leq \ (r_i)^{(3)} \leq (b_i^{'})^{(3)} \leq (\hat{B}_{20})^{(3)}$ | 112 |
| $lim_{T_2 \to \infty}(a_i^{''})^{(3)}\ (T_{21}, t) = (p_i)^{(3)}$ | 113 |

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|---|---|
| $\lim_{G\to\infty}(b_i'')^{(3)}(G_{23},t) = (r_i)^{(3)}$<br><br>**Definition of** $(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}$ :<br><br>Where $\boxed{(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}}$ are positive constants and $\boxed{i = 20,21,22}$ | |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i'')^{(3)}(T_{21}',t) - (a_i'')^{(3)}(T_{21},t)\rvert \le (\hat{k}_{20})^{(3)}\lvert T_{21} - T_{21}'\rvert e^{-(\hat{M}_{20})^{(3)}t}$<br><br>$\lvert(b_i'')^{(3)}(G_{23}',t) - (b_i'')^{(3)}(G_{23},t)\rvert < (\hat{k}_{20})^{(3)}\lvert\lvert G_{23} - G_{23}'\rvert\rvert e^{-(\hat{M}_{20})^{(3)}t}$ | 114 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(3)}(T_{21}',t)$ and$(a_i'')^{(3)}(T_{21},t)$ . $(T_{21}',t)$ And $(T_{21},t)$ are points belonging to the interval $\left[(\hat{k}_{20})^{(3)}, (\hat{M}_{20})^{(3)}\right]$. It is to be noted that $(a_i'')^{(3)}(T_{21},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{20})^{(3)} = 1$ then the function $(a_i'')^{(3)}(T_{21},t)$, the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$ :<br><br>$(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$, are positive constants<br>$\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}$ , $\frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}} < 1$ | 115 |
| There exists two constants There exists two constants $(\hat{P}_{20})^{(3)}$ and $(\hat{Q}_{20})^{(3)}$ which together with $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}, (\hat{A}_{20})^{(3)} and\ (\hat{B}_{20})^{(3)}$ and the constants $(a_i)^{(3)}, (a_i')^{(3)}, (b_i)^{(3)}, (b_i')^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}, i = 20,21,22,$ satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(a_i)^{(3)} + (a_i')^{(3)} + (\hat{A}_{20})^{(3)} + (\hat{P}_{20})^{(3)}(\hat{k}_{20})^{(3)}] < 1$<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(b_i)^{(3)} + (b_i')^{(3)} + (\hat{B}_{20})^{(3)} + (\hat{Q}_{20})^{(3)}(\hat{k}_{20})^{(3)}] < 1$ | 116 |
| Where we suppose | |
| $(a_i)^{(4)}, (a_i')^{(4)}, (a_i'')^{(4)}, (b_i)^{(4)}, (b_i')^{(4)}, (b_i'')^{(4)} > 0, \quad i,j = 24,25,26$<br><br>The functions $(a_i'')^{(4)}, (b_i'')^{(4)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(4)}$, $(r_i)^{(4)}$:<br><br>$(a_i'')^{(4)}(T_{25},t) \le (p_i)^{(4)} \le (\hat{A}_{24})^{(4)}$<br><br>$(b_i'')^{(4)}\big((G_{27}),t\big) \le (r_i)^{(4)} \le (b_i')^{(4)} \le (\hat{B}_{24})^{(4)}$ | 117 |
| $lim_{T_2\to\infty}(a_i'')^{(4)}(T_{25},t) = (p_i)^{(4)}$ | 118 |

| | |
|---|---|
| $\lim_{G\to\infty}(b_i'')^{(4)}\left((G_{27}),t\right) = \ (r_i)^{(4)}$<br><br>**Definition of** $(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}$ :<br><br>Where $\boxed{(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}, (p_i)^{(4)}, \ (r_i)^{(4)}}$ are positive constants and $\boxed{i = 24{,}25{,}26}$ | |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i'')^{(4)}(T_{25}',t) - (a_i'')^{(4)}(T_{25},t)\rvert \le (\hat{k}_{24})^{(4)}\lvert T_{25} - T_{25}'\rvert e^{-(\hat{M}_{24})^{(4)}t}$<br><br>$\lvert(b_i'')^{(4)}((G_{27})',t) - (b_i'')^{(4)}\left((G_{27}),t\right)\rvert < (\hat{k}_{24})^{(4)}\lvert\lvert(G_{27}) - (G_{27})'\rvert\rvert e^{-(\hat{M}_{24})^{(4)}t}$ | 119 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(4)}(T_{25}',t)$ and $(a_i'')^{(4)}(T_{25},t)$ . $(T_{25}',t)$ and $(T_{25},t)$ are points belonging to the interval $\left[(\hat{k}_{24})^{(4)}, (\hat{M}_{24})^{(4)}\right]$. It is to be noted that $(a_i'')^{(4)}(T_{25},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{24})^{(4)} = 1$ then the function $(a_i'')^{(4)}(T_{25},t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$ :<br><br>$(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$, are positive constants<br><br>$\frac{(a_i)^{(4)}}{(\hat{M}_{24})^{(4)}}$ , $\frac{(b_i)^{(4)}}{(\hat{M}_{24})^{(4)}} < 1$ | 120 |
| **Definition of** $(\hat{P}_{24})^{(4)}, (\hat{Q}_{24})^{(4)}$ :<br><br>There exists two constants $(\hat{P}_{24})^{(4)}$ and $(\hat{Q}_{24})^{(4)}$ which together with $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}, (\hat{A}_{24})^{(4)} and\ (\hat{B}_{24})^{(4)}$ and the constants $(a_i)^{(4)}, (a_i')^{(4)}, (b_i)^{(4)}, (b_i')^{(4)}, (p_i)^{(4)}, \ (r_i)^{(4)}, i = 24{,}25{,}26$, satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{24})^{(4)}}[\,(a_i)^{(4)} + (a_i')^{(4)} + \ (\hat{A}_{24})^{(4)} + \ (\hat{P}_{24})^{(4)}\,(\hat{k}_{24})^{(4)}] < 1$<br><br>$\frac{1}{(\hat{M}_{24})^{(4)}}[\ (b_i)^{(4)} + (b_i')^{(4)} + \ (\hat{B}_{24})^{(4)} + \ (\hat{Q}_{24})^{(4)}\ (\hat{k}_{24})^{(4)}] < 1$ | 121 |
| Where we suppose | |
| $(a_i)^{(5)}, (a_i')^{(5)}, (a_i'')^{(5)}, (b_i)^{(5)}, (b_i')^{(5)}, (b_i'')^{(5)} > 0, \quad i,j = 28{,}29{,}30$<br><br>The functions $(a_i'')^{(5)}, (b_i'')^{(5)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(5)}, \ (r_i)^{(5)}$:<br><br>$(a_i'')^{(5)}(T_{29},t) \le (p_i)^{(5)} \le (\hat{A}_{28})^{(5)}$<br><br>$(b_i'')^{(5)}\left((G_{31}),t\right) \le \ (r_i)^{(5)} \le (b_i')^{(5)} \le (\hat{B}_{28})^{(5)}$ | 122 |

$lim_{T_2\to\infty}(a_i'')^{(5)}(T_{29},t) = (p_i)^{(5)}$

$\lim_{G\to\infty}(b_i'')^{(5)}(G_{31},t) = (r_i)^{(5)}$

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**Definition of** $(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}$ :

Where $\boxed{(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}, (p_i)^{(5)}, (r_i)^{(5)}}$ are positive constants and $\boxed{i = 28,29,30}$

They satisfy Lipschitz condition:

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$$|(a_i'')^{(5)}(T_{29}', t) - (a_i'')^{(5)}(T_{29}, t)| \le (\hat{k}_{28})^{(5)}|T_{29} - T_{29}'|e^{-(\hat{M}_{28})^{(5)}t}$$

$$|(b_i'')^{(5)}((G_{31})', t) - (b_i'')^{(5)}((G_{31}), t)| < (\hat{k}_{28})^{(5)}||(G_{31}) - (G_{31})'||e^{-(\hat{M}_{28})^{(5)}t}$$

With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(5)}(T_{29}', t)$ and $(a_i'')^{(5)}(T_{29}, t)$ . $(T_{29}', t)$ and $(T_{29}, t)$ are points belonging to the interval $[(\hat{k}_{28})^{(5)}, (\hat{M}_{28})^{(5)}]$. It is to be noted that $(a_i'')^{(5)}(T_{29}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{28})^{(5)} = 1$ then the function $(a_i'')^{(5)}(T_{29}, t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

**Definition of** $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}$ :

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$(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)},$ are positive constants

$$\frac{(a_i)^{(5)}}{(\hat{M}_{28})^{(5)}}, \frac{(b_i)^{(5)}}{(\hat{M}_{28})^{(5)}} < 1$$

**Definition of** $(\hat{P}_{28})^{(5)}, (\hat{Q}_{28})^{(5)}$ :

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There exists two constants $(\hat{P}_{28})^{(5)}$ and $(\hat{Q}_{28})^{(5)}$ which together with $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}, (\hat{A}_{28})^{(5)} and\ (\hat{B}_{28})^{(5)}$ and the constants $(a_i)^{(5)}, (a_i')^{(5)}, (b_i)^{(5)}, (b_i')^{(5)}, (p_i)^{(5)}, (r_i)^{(5)}, i = 28,29,30,$ satisfy the inequalities

$$\frac{1}{(\hat{M}_{28})^{(5)}}[(a_i)^{(5)} + (a_i')^{(5)} + (\hat{A}_{28})^{(5)} + (\hat{P}_{28})^{(5)}(\hat{k}_{28})^{(5)}] < 1$$

$$\frac{1}{(\hat{M}_{28})^{(5)}}[(b_i)^{(5)} + (b_i')^{(5)} + (\hat{B}_{28})^{(5)} + (\hat{Q}_{28})^{(5)}(\hat{k}_{28})^{(5)}] < 1$$

Where we suppose

$(a_i)^{(6)}, (a_i')^{(6)}, (a_i'')^{(6)}, (b_i)^{(6)}, (b_i')^{(6)}, (b_i'')^{(6)} > 0, \quad i, j = 32,33,34$

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The functions $(a_i'')^{(6)}, (b_i'')^{(6)}$ are positive continuous increasing and bounded.

**Definition of** $(p_i)^{(6)}, (r_i)^{(6)}$:

$$(a_i'')^{(6)}(T_{33}, t) \le (p_i)^{(6)} \le (\hat{A}_{32})^{(6)}$$

| | |
|---|---|
| $(b_i^{''})^{(6)}((G_{35}),t) \leq \ (r_i)^{(6)} \leq (b_i^{'})^{(6)} \leq (\hat{B}_{32})^{(6)}$ | |
| $lim_{T_2\to\infty}(a_i^{''})^{(6)}(T_{33},t) = (p_i)^{(6)}$<br><br>$\lim_{G\to\infty}(b_i^{''})^{(6)}\big((G_{35}),t\big) = \ (r_i)^{(6)}$<br><br>**Definition of** $(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}$ :<br><br>Where $\boxed{(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}, (p_i)^{(6)}, \ (r_i)^{(6)}}$ are positive constants and $\boxed{i = 32,33,34}$ | 128 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(6)}(T_{33}^{'},t) - (a_i^{''})^{(6)}(T_{33},t)\rvert \leq (\hat{k}_{32})^{(6)}\lvert T_{33} - T_{33}^{'}\rvert e^{-(\hat{M}_{32})^{(6)}t}$<br><br>$\lvert(b_i^{''})^{(6)}((G_{35})^{'},t) - (b_i^{''})^{(6)}\big((G_{35}),t\big)\rvert < (\hat{k}_{32})^{(6)}\lvert\lvert(G_{35}) - (G_{35})^{'}\rvert\rvert e^{-(\hat{M}_{32})^{(6)}t}$ | |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(6)}(T_{33}^{'},t)$ and$(a_i^{''})^{(6)}(T_{33},t)$ . $(T_{33}^{'},t)$ and $(T_{33},t)$ are points belonging to the interval $\big[(\hat{k}_{32})^{(6)}, (\hat{M}_{32})^{(6)}\big]$. It is to be noted that $(a_i^{''})^{(6)}(T_{33},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{32})^{(6)} = 1$ then the function $(a_i^{''})^{(6)}(T_{33},t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}$ :<br><br>$(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}$, are positive constants<br><br>$\frac{(a_i)^{(6)}}{(\hat{M}_{32})^{(6)}}$ , $\frac{(b_i)^{(6)}}{(\hat{M}_{32})^{(6)}} < 1$ | 129 |
| **Definition of** $(\hat{P}_{32})^{(6)}, (\hat{Q}_{32})^{(6)}$ :<br><br>There exists two constants $(\hat{P}_{32})^{(6)}$ and $(\hat{Q}_{32})^{(6)}$ which together with $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}, (\hat{A}_{32})^{(6)} and\ (\hat{B}_{32})^{(6)}$ and the constants $(a_i)^{(6)}, (a_i^{'})^{(6)}, (b_i)^{(6)}, (b_i^{'})^{(6)}, (p_i)^{(6)}, \ (r_i)^{(6)}, i = 32,33,34,$ satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{32})^{(6)}}[\,(a_i)^{(6)} + (a_i^{'})^{(6)} + \ (\hat{A}_{32})^{(6)} + \ (\hat{P}_{32})^{(6)}\,(\hat{k}_{32})^{(6)}] < 1$<br><br>$\frac{1}{(\hat{M}_{32})^{(6)}}[\ (b_i)^{(6)} + (b_i^{'})^{(6)} + \ (\hat{B}_{32})^{(6)} + \ (\hat{Q}_{32})^{(6)}\ (\hat{k}_{32})^{(6)}] < 1$ | 130 |
| Where we suppose | |
| (S) $(a_i)^{(7)}, (a_i^{'})^{(7)}, (a_i^{''})^{(7)}, (b_i)^{(7)}, (b_i^{'})^{(7)}, (b_i^{''})^{(7)} > 0, \quad i,j = 36,37,38$<br><br>(T) The functions $(a_i^{''})^{(7)}, (b_i^{''})^{(7)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(7)}, \ (r_i)^{(7)}$: | 131 |

$(a_i^{''})^{(7)}(T_{37},t) \leq (p_i)^{(7)} \leq (\hat{A}_{36})^{(7)}$

$(b_i^{''})^{(7)}(G_{39},t) \leq (r_i)^{(7)} \leq (b_i^{'})^{(7)} \leq (\hat{B}_{36})^{(7)}$

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*(U)* $\lim_{T_2\to\infty} (a_i^{''})^{(7)}(T_{37},t) = (p_i)^{(7)}$

*(V)*

$\lim_{G\to\infty}(b_i^{''})^{(7)}\left((G_{39}),t\right) = (r_i)^{(7)}$

**Definition of** $(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}$ :

Where $\boxed{(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}, (p_i)^{(7)}, (r_i)^{(7)}}$ are positive constants and $\boxed{i = 36,37,38}$

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They satisfy Lipschitz condition:

$|(a_i^{''})^{(7)}(T_{37}^{'},t) - (a_i^{''})^{(7)}(T_{37},t)| \leq (\hat{k}_{36})^{(7)}|T_{37} - T_{37}^{'}|e^{-(\hat{M}_{36})^{(7)}t}$

$|(b_i^{''})^{(7)}((G_{39})^{'},t) - (b_i^{''})^{(7)}\left((G_{39}),t\right)| < (\hat{k}_{36})^{(7)}||(G_{39}) - (G_{39})^{'}||e^{-(\hat{M}_{36})^{(7)}t}$

With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(7)}(T_{37}^{'},t)$ and$(a_i^{''})^{(7)}(T_{37},t)$ . $(T_{37}^{'},t)$ and $(T_{37},t)$ are points belonging to the interval $\left[(\hat{k}_{36})^{(7)}, (\hat{M}_{36})^{(7)}\right]$. It is to be noted that $(a_i^{''})^{(7)}(T_{37},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{36})^{(7)} = 1$ then the function $(a_i^{''})^{(7)}(T_{37},t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

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**Definition of** $(\hat{M}_{36})^{(7)}, (\hat{k}_{36})^{(7)}$ :

(W) $(\hat{M}_{36})^{(7)}, (\hat{k}_{36})^{(7)}$, are positive constants

$\frac{(a_i)^{(7)}}{(\hat{M}_{36})^{(7)}}, \frac{(b_i)^{(7)}}{(\hat{M}_{36})^{(7)}} < 1$

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**Definition of** $(\hat{P}_{36})^{(7)}, (\hat{Q}_{36})^{(7)}$ :

(X) There exists two constants $(\hat{P}_{36})^{(7)}$ and $(\hat{Q}_{36})^{(7)}$ which together with $(\hat{M}_{36})^{(7)}, (\hat{k}_{36})^{(7)}, (\hat{A}_{36})^{(7)} and\ (\hat{B}_{36})^{(7)}$ and the constants $(a_i)^{(7)}, (a_i^{'})^{(7)}, (b_i)^{(7)}, (b_i^{'})^{(7)}, (p_i)^{(7)}, (r_i)^{(7)}, i = 36,37,38$, satisfy the inequalities

$\frac{1}{(\hat{M}_{36})^{(7)}}[(a_i)^{(7)} + (a_i^{'})^{(7)} + (\hat{A}_{36})^{(7)} + (\hat{P}_{36})^{(7)}(\hat{k}_{36})^{(7)}] < 1$

$\frac{1}{(\hat{M}_{36})^{(7)}}[(b_i)^{(7)} + (b_i^{'})^{(7)} + (\hat{B}_{36})^{(7)} + (\hat{Q}_{36})^{(7)}(\hat{k}_{36})^{(7)}] < 1$

| | |
|---|---|
| Where we suppose | |
| $(a_i)^{(8)}, (a_i^{'})^{(8)}, (a_i^{''})^{(8)}, (b_i)^{(8)}, (b_i^{'})^{(8)}, (b_i^{''})^{(8)} > 0, \quad i,j = 40,41,42$ | 136 |
| The functions $(a_i^{''})^{(8)}, (b_i^{''})^{(8)}$ are positive continuous increasing and bounded | |
| **Definition of** $(p_i)^{(8)}$, $(r_i)^{(8)}$: | 137 |
| $(a_i^{''})^{(8)}(T_{41}, t) \leq (p_i)^{(8)} \leq (\hat{A}_{40})^{(8)}$ | 138 |
| $(b_i^{''})^{(8)}((G_{43}), t) \leq (r_i)^{(8)} \leq (b_i^{'})^{(8)} \leq (\hat{B}_{40})^{(8)}$ | 139 |
| $\lim_{T_2 \to \infty} (a_i^{''})^{(8)}(T_{41}, t) = (p_i)^{(8)}$ | 140 |
| $\lim_{G \to \infty} (b_i^{''})^{(8)}((G_{43}), t) = (r_i)^{(8)}$ | 141 |
| **Definition of** $(\hat{A}_{40})^{(8)}, (\hat{B}_{40})^{(8)}$ :<br>Where $\boxed{(\hat{A}_{40})^{(8)}, (\hat{B}_{40})^{(8)}, (p_i)^{(8)}, (r_i)^{(8)}}$ are positive constants and $\boxed{i = 40,41,42}$ | |
| They satisfy Lipschitz condition: | |
| $\lvert(a_i^{''})^{(8)}(T_{41}^{'}, t) - (a_i^{''})^{(8)}(T_{41}, t)\rvert \leq (\hat{k}_{40})^{(8)} \lvert T_{41} - T_{41}^{'}\rvert e^{-(\hat{M}_{40})^{(8)} t}$ | 142 |
| $\lvert(b_i^{''})^{(8)}((G_{43})^{'}, t) - (b_i^{''})^{(8)}((G_{43}), t)\rvert < (\hat{k}_{40})^{(8)} \lvert\lvert(G_{43}) - (G_{43})^{'}\rvert\rvert e^{-(\hat{M}_{40})^{(8)} t}$ | 143 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(8)}(T_{41}^{'}, t)$ and $(a_i^{''})^{(8)}(T_{41}, t)$ . $(T_{41}^{'}, t)$ and $(T_{41}, t)$ are points belonging to the interval $\left[(\hat{k}_{40})^{(8)}, (\hat{M}_{40})^{(8)}\right]$. It is to be noted that $(a_i^{''})^{(8)}(T_{41}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{40})^{(8)} = 1$ then the function $(a_i^{''})^{(8)}(T_{41}, t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}$ : | |
| $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}$, are positive constants | |
| $\frac{(a_i)^{(8)}}{(\hat{M}_{40})^{(8)}}$ , $\frac{(b_i)^{(8)}}{(\hat{M}_{40})^{(8)}} < 1$ | 144 |
| **Definition of** $(\hat{P}_{40})^{(8)}, (\hat{Q}_{40})^{(8)}$ :<br>There exists two constants $(\hat{P}_{40})^{(8)}$ and $(\hat{Q}_{40})^{(8)}$ which together with $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}, (\hat{A}_{40})^{(8)}$ $(\hat{B}_{40})^{(8)}$ and the constants $(a_i)^{(8)}, (a_i^{'})^{(8)}, (b_i)^{(8)}, (b_i^{'})^{(8)}, (p_i)^{(8)}$, $(r_i)^{(8)}, i = 40,41,42,$<br>Satisfy the inequalities | |

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| $\frac{1}{(\hat{M}_{40})^{(8)}}[(a_i)^{(8)} + (a_i')^{(8)} + (\hat{A}_{40})^{(8)} + (\hat{P}_{40})^{(8)}(\hat{k}_{40})^{(8)}] < 1$ | 145 |
| $\frac{1}{(\hat{M}_{40})^{(8)}}[(b_i)^{(8)} + (b_i')^{(8)} + (\hat{B}_{40})^{(8)} + (\hat{Q}_{40})^{(8)}(\hat{k}_{40})^{(8)}] < 1$ | 146 |
| Where we suppose | |
| $(a_i)^{(9)}, (a_i')^{(9)}, (a_i'')^{(9)}, (b_i)^{(9)}, (b_i')^{(9)}, (b_i'')^{(9)} > 0, \quad i,j = 44,45,46$<br>The functions $(a_i'')^{(9)}, (b_i'')^{(9)}$ are positive continuous increasing and bounded.<br>**Definition of** $(p_i)^{(9)}$, $(r_i)^{(9)}$:<br>$(a_i'')^{(9)}(T_{45},t) \le (p_i)^{(9)} \le (\hat{A}_{44})^{(9)}$<br>$(b_i'')^{(9)}(G_{47},t) \le (r_i)^{(9)} \le (b_i')^{(9)} \le (\hat{B}_{44})^{(9)}$ | 146 A |
| $\lim_{T_2\to\infty}(a_i'')^{(9)}(T_{45},t) = (p_i)^{(9)}$<br>$\lim_{G\to\infty}(b_i'')^{(9)}(G_{47},t) = (r_i)^{(9)}$<br>**Definition of** $(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}$ :<br>Where $\boxed{(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}, (p_i)^{(9)}, (r_i)^{(9)}}$ are positive constants and $\boxed{i = 44,45,46}$ | |
| They satisfy Lipschitz condition:<br>$\lvert(a_i'')^{(9)}(T_{45}',t) - (a_i'')^{(9)}(T_{45},t)\rvert \le (\hat{k}_{44})^{(9)}\lvert T_{45} - T_{45}'\rvert e^{-(\hat{M}_{44})^{(9)}t}$<br>$\lvert(b_i'')^{(9)}((G_{47})',t) - (b_i'')^{(9)}((G_{47}),t)\rvert < (\hat{k}_{44})^{(9)}\lvert\lvert(G_{47}) - (G_{47})'\rvert\rvert e^{-(\hat{M}_{44})^{(9)}t}$ | |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(9)}(T_{45}',t)$ and $(a_i'')^{(9)}(T_{45},t)$ . $(T_{45}',t)$ and $(T_{45},t)$ are points belonging to the interval $[(\hat{k}_{44})^{(9)}, (\hat{M}_{44})^{(9)}]$. It is to be noted that $(a_i'')^{(9)}(T_{45},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{44})^{(9)} = 1$ then the function $(a_i'')^{(9)}(T_{45},t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}$ :<br>$(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}$, are positive constants<br>$\frac{(a_i)^{(9)}}{(\hat{M}_{44})^{(9)}}, \frac{(b_i)^{(9)}}{(\hat{M}_{44})^{(9)}} < 1$ | |
| **Definition of** $(\hat{P}_{44})^{(9)}, (\hat{Q}_{44})^{(9)}$ :<br>There exists two constants $(\hat{P}_{44})^{(9)}$ and $(\hat{Q}_{44})^{(9)}$ which together with $(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}, (\hat{A}_{44})^{(9)} and\ (\hat{B}_{44})^{(9)}$ and the constants | |

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| $(a_i)^{(9)}, (a_i')^{(9)}, (b_i)^{(9)}, (b_i')^{(9)}, (p_i)^{(9)}, \ (r_i)^{(9)}, i = 44,45,46,$ satisfy the inequalities<br>$\frac{1}{(\hat{M}_{44})^{(9)}}[(a_i)^{(9)} + (a_i')^{(9)} + (\hat{A}_{44})^{(9)} + (\hat{P}_{44})^{(9)} (\hat{k}_{44})^{(9)}] < 1$<br>$\frac{1}{(\hat{M}_{44})^{(9)}}[(b_i)^{(9)} + (b_i')^{(9)} + (\hat{B}_{44})^{(9)} + (\hat{Q}_{44})^{(9)} (\hat{k}_{44})^{(9)}] < 1$ | |
| **Theorem 1:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>Definition of $G_i(0), T_i(0)$ :<br>$G_i(t) \le (\hat{P}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le (\hat{Q}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 147 |
| **Theorem 2 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$<br>$G_i(t) \le (\hat{P}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ , $G_i(0) = G_i^0 > 0$<br>$T_i(t) \le (\hat{Q}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ , $T_i(0) = T_i^0 > 0$ | 148 |
| **Theorem 3 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>$G_i(t) \le (\hat{P}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ , $G_i(0) = G_i^0 > 0$<br>$T_i(t) \le (\hat{Q}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ , $T_i(0) = T_i^0 > 0$ | 149 |
| **Theorem 4 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \le (\hat{P}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le (\hat{Q}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 150 |
| **Theorem 5 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \le (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le (\hat{Q}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 151 |

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| **Theorem 6 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br><br>**Definition of** $G_i(0)\,, T_i(0)$ :<br><br>$G_i(t) \le \left(\hat{P}_{32}\right)^{(6)} e^{(\hat{M}_{32})^{(6)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br><br>$T_i(t) \le \left(\hat{Q}_{32}\right)^{(6)} e^{(\hat{M}_{32})^{(6)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 152 |
| **Theorem 7:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br><br>**Definition of** $G_i(0)\,, T_i(0)$ :<br><br>$G_i(t) \le \left(\hat{P}_{36}\right)^{(7)} e^{(\hat{M}_{36})^{(7)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{36}\right)^{(7)} e^{(\hat{M}_{36})^{(7)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 |
| **Theorem 8:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br><br>**Definition of** $G_i(0)\,, T_i(0)$ :<br><br>$G_i(t) \le \left(\hat{P}_{40}\right)^{(8)} e^{(\hat{M}_{40})^{(8)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{40}\right)^{(8)} e^{(\hat{M}_{40})^{(8)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 A |
| **Theorem 9:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br><br>**Definition of** $G_i(0)\,, T_i(0)$ :<br><br>$G_i(t) \le \left(\hat{P}_{44}\right)^{(9)} e^{(\hat{M}_{44})^{(9)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br><br>$T_i(t) \le \left(\hat{Q}_{44}\right)^{(9)} e^{(\hat{M}_{44})^{(9)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 B |
| **Proof:** Consider operator $\mathcal{A}^{(1)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i$: $\mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 154 |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \le \left(\hat{P}_{13}\right)^{(1)}$ , $T_i^0 \le \left(\hat{Q}_{13}\right)^{(1)}$, | 155 |
| $0 \le G_i(t) - G_i^0 \le \left(\hat{P}_{13}\right)^{(1)} e^{(\hat{M}_{13})^{(1)}t}$ | 156 |
| $0 \le T_i(t) - T_i^0 \le \left(\hat{Q}_{13}\right)^{(1)} e^{(\hat{M}_{13})^{(1)}t}$ | 157 |
| By<br><br>$\bar{G}_{13}(t) = G_{13}^0 + \int_0^t \left[(a_{13})^{(1)} G_{14}\left(s_{(13)}\right) - \left((a_{13}')^{(1)} + a_{13}'')^{(1)}\left(T_{14}\left(s_{(13)}\right), s_{(13)}\right)\right) G_{13}\left(s_{(13)}\right)\right] ds_{(13)}$ | 158 |

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| $\bar{G}_{14}(t) = G_{14}^0 + \int_0^t \left[(a_{14})^{(1)} G_{13}(s_{(13)}) - \left((a_{14}')^{(1)} + (a_{14}'')^{(1)}(T_{14}(s_{(13)}), s_{(13)})\right) G_{14}(s_{(13)})\right] ds_{(13)}$ | |
| $\bar{G}_{15}(t) = G_{15}^0 + \int_0^t \left[(a_{15})^{(1)} G_{14}(s_{(13)}) - \left((a_{15}')^{(1)} + (a_{15}'')^{(1)}(T_{14}(s_{(13)}), s_{(13)})\right) G_{15}(s_{(13)})\right] ds_{(13)}$ | |
| $\bar{T}_{13}(t) = T_{13}^0 + \int_0^t \left[(b_{13})^{(1)} T_{14}(s_{(13)}) - \left((b_{13}')^{(1)} - (b_{13}'')^{(1)}(G(s_{(13)}), s_{(13)})\right) T_{13}(s_{(13)})\right] ds_{(13)}$ | |
| $\bar{T}_{14}(t) = T_{14}^0 + \int_0^t \left[(b_{14})^{(1)} T_{13}(s_{(13)}) - \left((b_{14}')^{(1)} - (b_{14}'')^{(1)}(G(s_{(13)}), s_{(13)})\right) T_{14}(s_{(13)})\right] ds_{(13)}$ | |
| $\bar{T}_{15}(t) = T_{15}^0 + \int_0^t \left[(b_{15})^{(1)} T_{14}(s_{(13)}) - \left((b_{15}')^{(1)} - (b_{15}'')^{(1)}(G(s_{(13)}), s_{(13)})\right) T_{15}(s_{(13)})\right] ds_{(13)}$<br>Where $s_{(13)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(2)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 159 |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{16})^{(2)}$, $T_i^0 \le (\hat{Q}_{16})^{(2)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ | |
| By<br>$\bar{G}_{16}(t) = G_{16}^0 + \int_0^t \left[(a_{16})^{(2)} G_{17}(s_{(16)}) - \left((a_{16}')^{(2)} + a_{16}'')^{(2)}(T_{17}(s_{(16)}), s_{(16)})\right) G_{16}(s_{(16)})\right] ds_{(16)}$ | 160 |
| $\bar{G}_{17}(t) = G_{17}^0 + \int_0^t \left[(a_{17})^{(2)} G_{16}(s_{(16)}) - \left((a_{17}')^{(2)} + (a_{17}'')^{(2)}(T_{17}(s_{(16)}), s_{(17)})\right) G_{17}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{G}_{18}(t) = G_{18}^0 + \int_0^t \left[(a_{18})^{(2)} G_{17}(s_{(16)}) - \left((a_{18}')^{(2)} + (a_{18}'')^{(2)}(T_{17}(s_{(16)}), s_{(16)})\right) G_{18}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{T}_{16}(t) = T_{16}^0 + \int_0^t \left[(b_{16})^{(2)} T_{17}(s_{(16)}) - \left((b_{16}')^{(2)} - (b_{16}'')^{(2)}(G_{19}(s_{(16)}), s_{(16)})\right) T_{16}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{T}_{17}(t) = T_{17}^0 + \int_0^t \left[(b_{17})^{(2)} T_{16}(s_{(16)}) - \left((b_{17}')^{(2)} - (b_{17}'')^{(2)}(G_{19}(s_{(16)}), s_{(16)})\right) T_{17}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{T}_{18}(t) = T_{18}^0 + \int_0^t \left[(b_{18})^{(2)} T_{17}(s_{(16)}) - \left((b_{18}')^{(2)} - (b_{18}'')^{(2)}(G_{19}(s_{(16)}), s_{(16)})\right) T_{18}(s_{(16)})\right] ds_{(16)}$<br>Where $s_{(16)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(3)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |

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| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{20})^{(3)}$ , $T_i^0 \leq (\hat{Q}_{20})^{(3)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{20})^{(3)} e^{(\bar{M}_{20})^{(3)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{20})^{(3)} e^{(\bar{M}_{20})^{(3)} t}$ | |
| By<br>$\bar{G}_{20}(t) = G_{20}^0 + \int_0^t \left[ (a_{20})^{(3)} G_{21}(s_{(20)}) - \left( (a'_{20})^{(3)} + a''_{20})^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{20}(s_{(20)}) \right] ds_{(20)}$ | 161 |
| $\bar{G}_{21}(t) = G_{21}^0 + \int_0^t \left[ (a_{21})^{(3)} G_{20}(s_{(20)}) - \left( (a'_{21})^{(3)} + (a''_{21})^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{21}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{G}_{22}(t) = G_{22}^0 + \int_0^t \left[ (a_{22})^{(3)} G_{21}(s_{(20)}) - \left( (a'_{22})^{(3)} + (a''_{22})^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{22}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{20}(t) = T_{20}^0 + \int_0^t \left[ (b_{20})^{(3)} T_{21}(s_{(20)}) - \left( (b'_{20})^{(3)} - (b''_{20})^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{20}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{21}(t) = T_{21}^0 + \int_0^t \left[ (b_{21})^{(3)} T_{20}(s_{(20)}) - \left( (b'_{21})^{(3)} - (b''_{21})^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{21}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{22}(t) = T_{22}^0 + \int_0^t \left[ (b_{22})^{(3)} T_{21}(s_{(20)}) - \left( (b'_{22})^{(3)} - (b''_{22})^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{22}(s_{(20)}) \right] ds_{(20)}$<br>Where $s_{(20)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:** Consider operator $\mathcal{A}^{(4)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i$: $\mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{24})^{(4)}$ , $T_i^0 \leq (\hat{Q}_{24})^{(4)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{24})^{(4)} e^{(\bar{M}_{24})^{(4)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{24})^{(4)} e^{(\bar{M}_{24})^{(4)} t}$ | |
| By<br>$\bar{G}_{24}(t) = G_{24}^0 + \int_0^t \left[ (a_{24})^{(4)} G_{25}(s_{(24)}) - \left( (a'_{24})^{(4)} + a''_{24})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{24}(s_{(24)}) \right] ds_{(24)}$ | 162 |
| $\bar{G}_{25}(t) = G_{25}^0 + \int_0^t \left[ (a_{25})^{(4)} G_{24}(s_{(24)}) - \left( (a'_{25})^{(4)} + (a''_{25})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{25}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{G}_{26}(t) = G_{26}^0 + \int_0^t \left[ (a_{26})^{(4)} G_{25}(s_{(24)}) - \left( (a'_{26})^{(4)} + (a''_{26})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{26}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{24}(t) = T_{24}^0 + \int_0^t \left[ (b_{24})^{(4)} T_{25}(s_{(24)}) - \left( (b'_{24})^{(4)} - (b''_{24})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{24}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{25}(t) = T_{25}^0 + \int_0^t \left[ (b_{25})^{(4)} T_{24}(s_{(24)}) - \left( (b'_{25})^{(4)} - (b''_{25})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{25}(s_{(24)}) \right] ds_{(24)}$ | |

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| $\overline{T}_{26}(t) = T_{26}^0 + \int_0^t \left[(b_{26})^{(4)} T_{25}(s_{(24)}) - \left((b_{26}')^{(4)} - (b_{26}'')^{(4)}(G_{27}(s_{(24)}), s_{(24)})\right) T_{26}(s_{(24)})\right] ds_{(24)}$<br>Where $s_{(24)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:** Consider operator $\mathcal{A}^{(5)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{28})^{(5)}$ , $T_i^0 \leq (\hat{Q}_{28})^{(5)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| By<br>$\bar{G}_{28}(t) = G_{28}^0 + \int_0^t \left[(a_{28})^{(5)} G_{29}(s_{(28)}) - \left((a_{28}')^{(5)} + a_{28}'')^{(5)}(T_{29}(s_{(28)}), s_{(28)})\right) G_{28}(s_{(28)})\right] ds_{(28)}$ | 163 |
| $\bar{G}_{29}(t) = G_{29}^0 + \int_0^t \left[(a_{29})^{(5)} G_{28}(s_{(28)}) - \left((a_{29}')^{(5)} + (a_{29}'')^{(5)}(T_{29}(s_{(28)}), s_{(28)})\right) G_{29}(s_{(28)})\right] ds_{(28)}$ | |
| $\bar{G}_{30}(t) = G_{30}^0 + \int_0^t \left[(a_{30})^{(5)} G_{29}(s_{(28)}) - \left((a_{30}')^{(5)} + (a_{30}'')^{(5)}(T_{29}(s_{(28)}), s_{(28)})\right) G_{30}(s_{(28)})\right] ds_{(28)}$ | |
| $\overline{T}_{28}(t) = T_{28}^0 + \int_0^t \left[(b_{28})^{(5)} T_{29}(s_{(28)}) - \left((b_{28}')^{(5)} - (b_{28}'')^{(5)}(G_{31}(s_{(28)}), s_{(28)})\right) T_{28}(s_{(28)})\right] ds_{(28)}$ | |
| $\overline{T}_{29}(t) = T_{29}^0 + \int_0^t \left[(b_{29})^{(5)} T_{28}(s_{(28)}) - \left((b_{29}')^{(5)} - (b_{29}'')^{(5)}(G_{31}(s_{(28)}), s_{(28)})\right) T_{29}(s_{(28)})\right] ds_{(28)}$ | |
| $\overline{T}_{30}(t) = T_{30}^0 + \int_0^t \left[(b_{30})^{(5)} T_{29}(s_{(28)}) - \left((b_{30}')^{(5)} - (b_{30}'')^{(5)}(G_{31}(s_{(28)}), s_{(28)})\right) T_{30}(s_{(28)})\right] ds_{(28)}$<br>Where $s_{(28)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(6)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{32})^{(6)}$ , $T_i^0 \leq (\hat{Q}_{32})^{(6)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ | |
| By<br>$\bar{G}_{32}(t) = G_{32}^0 + \int_0^t \left[(a_{32})^{(6)} G_{33}(s_{(32)}) - \left((a_{32}')^{(6)} + a_{32}'')^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{32}(s_{(32)})\right] ds_{(32)}$ | 164 |
| $\bar{G}_{33}(t) = G_{33}^0 + \int_0^t \left[(a_{33})^{(6)} G_{32}(s_{(32)}) - \left((a_{33}')^{(6)} + (a_{33}'')^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{33}(s_{(32)})\right] ds_{(32)}$ | |

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| $\bar{G}_{34}(t) = G_{34}^0 + \int_0^t \left[(a_{34})^{(6)} G_{33}(s_{(32)}) - \left((a_{34}')^{(6)} + (a_{34}'')^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{34}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{T}_{32}(t) = T_{32}^0 + \int_0^t \left[(b_{32})^{(6)} T_{33}(s_{(32)}) - \left((b_{32}')^{(6)} - (b_{32}'')^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{32}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{T}_{33}(t) = T_{33}^0 + \int_0^t \left[(b_{33})^{(6)} T_{32}(s_{(32)}) - \left((b_{33}')^{(6)} - (b_{33}'')^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{33}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{T}_{34}(t) = T_{34}^0 + \int_0^t \left[(b_{34})^{(6)} T_{33}(s_{(32)}) - \left((b_{34}')^{(6)} - (b_{34}'')^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{34}(s_{(32)})\right] ds_{(32)}$<br>Where $s_{(32)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(7)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{36})^{(7)}$, $T_i^0 \le (\hat{Q}_{36})^{(7)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ | |
| By<br>$\bar{G}_{36}(t) = G_{36}^0 + \int_0^t \left[(a_{36})^{(7)} G_{37}(s_{(36)}) - \left((a_{36}')^{(7)} + a_{36}'')^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{36}(s_{(36)})\right] ds_{(36)}$ | 165 |
| $\bar{G}_{37}(t) = G_{37}^0 + \int_0^t \left[(a_{37})^{(7)} G_{36}(s_{(36)}) - \left((a_{37}')^{(7)} + (a_{37}'')^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{37}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{G}_{38}(t) = G_{38}^0 + \int_0^t \left[(a_{38})^{(7)} G_{37}(s_{(36)}) - \left((a_{38}')^{(7)} + (a_{38}'')^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{38}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{36}(t) = T_{36}^0 + \int_0^t \left[(b_{36})^{(7)} T_{37}(s_{(36)}) - \left((b_{36}')^{(7)} - (b_{36}'')^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{36}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{37}(t) = T_{37}^0 + \int_0^t \left[(b_{37})^{(7)} T_{36}(s_{(36)}) - \left((b_{37}')^{(7)} - (b_{37}'')^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{37}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{38}(t) = T_{38}^0 + \int_0^t \left[(b_{38})^{(7)} T_{37}(s_{(36)}) - \left((b_{38}')^{(7)} - (b_{38}'')^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{38}(s_{(36)})\right] ds_{(36)}$<br>Where $s_{(36)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(8)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ | |

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| which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{40})^{(8)}$ , $T_i^0 \leq (\hat{Q}_{40})^{(8)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{40})^{(8)} e^{(\bar{M}_{40})^{(8)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{40})^{(8)} e^{(\bar{M}_{40})^{(8)} t}$ | |
| By<br>$\bar{G}_{40}(t) = G_{40}^0 + \int_0^t \left[ (a_{40})^{(8)} G_{41}(s_{(40)}) - \left( (a'_{40})^{(8)} + a''_{40})^{(8)} (T_{41}(s_{(40)}), s_{(40)}) \right) G_{40}(s_{(40)}) \right] ds_{(40)}$ | 166 |
| $\bar{G}_{41}(t) = G_{41}^0 + \int_0^t \left[ (a_{41})^{(8)} G_{40}(s_{(40)}) - \left( (a'_{41})^{(8)} + (a''_{41})^{(8)} (T_{41}(s_{(40)}), s_{(40)}) \right) G_{41}(s_{(40)}) \right] ds_{(40)}$ | |
| $\bar{G}_{42}(t) = G_{42}^0 + \int_0^t \left[ (a_{42})^{(8)} G_{41}(s_{(40)}) - \left( (a'_{42})^{(8)} + (a''_{42})^{(8)} (T_{41}(s_{(40)}), s_{(40)}) \right) G_{42}(s_{(40)}) \right] ds_{(40)}$ | |
| $\bar{T}_{40}(t) = T_{40}^0 + \int_0^t \left[ (b_{40})^{(8)} T_{41}(s_{(40)}) - \left( (b'_{40})^{(8)} - (b''_{40})^{(8)} (G_{43}(s_{(40)}), s_{(40)}) \right) T_{40}(s_{(40)}) \right] ds_{(40)}$ | |
| $\bar{T}_{41}(t) = T_{41}^0 + \int_0^t \left[ (b_{41})^{(8)} T_{40}(s_{(40)}) - \left( (b'_{41})^{(8)} - (b''_{41})^{(8)} (G_{43}(s_{(40)}), s_{(40)}) \right) T_{41}(s_{(40)}) \right] ds_{(40)}$ | |
| $\bar{T}_{42}(t) = T_{42}^0 + \int_0^t \left[ (b_{42})^{(8)} T_{41}(s_{(40)}) - \left( (b'_{42})^{(8)} - (b''_{42})^{(8)} (G_{43}(s_{(40)}), s_{(40)}) \right) T_{42}(s_{(40)}) \right] ds_{(40)}$<br>Where $s_{(40)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(9)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 166 A |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{44})^{(9)}$ , $T_i^0 \leq (\hat{Q}_{44})^{(9)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{44})^{(9)} e^{(\bar{M}_{44})^{(9)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{44})^{(9)} e^{(\bar{M}_{44})^{(9)} t}$ | |
| By<br>$\bar{G}_{44}(t) = G_{44}^0 + \int_0^t \left[ (a_{44})^{(9)} G_{45}(s_{(44)}) - \left( (a'_{44})^{(9)} + a''_{44})^{(9)} (T_{45}(s_{(44)}), s_{(44)}) \right) G_{44}(s_{(44)}) \right] ds_{(44)}$ | |
| $\bar{G}_{45}(t) = G_{45}^0 + \int_0^t \left[ (a_{45})^{(9)} G_{44}(s_{(44)}) - \left( (a'_{45})^{(9)} + (a''_{45})^{(9)} (T_{45}(s_{(44)}), s_{(44)}) \right) G_{45}(s_{(44)}) \right] ds_{(44)}$ | |
| $\bar{G}_{46}(t) = G_{46}^0 + \int_0^t \left[ (a_{46})^{(9)} G_{45}(s_{(44)}) - \left( (a'_{46})^{(9)} + (a''_{46})^{(9)} (T_{45}(s_{(44)}), s_{(44)}) \right) G_{46}(s_{(44)}) \right] ds_{(44)}$ | |

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| $\bar{T}_{44}(t) = T_{44}^0 + \int_0^t \left[(b_{44})^{(9)} T_{45}(s_{(44)}) - \left((b'_{44})^{(9)} - (b''_{44})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{44}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{45}(t) = T_{45}^0 + \int_0^t \left[(b_{45})^{(9)} T_{44}(s_{(44)}) - \left((b'_{45})^{(9)} - (b''_{45})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{45}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{46}(t) = T_{46}^0 + \int_0^t \left[(b_{46})^{(9)} T_{45}(s_{(44)}) - \left((b'_{46})^{(9)} - (b''_{46})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{46}(s_{(44)})\right] ds_{(44)}$<br>Where $s_{(44)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| The operator $\mathcal{A}^{(1)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{13}(t) \le G_{13}^0 + \int_0^t \left[(a_{13})^{(1)} \left(G_{14}^0 + (\hat{P}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} s_{(13)}}\right)\right] ds_{(13)} =$<br>$\left(1 + (a_{13})^{(1)} t\right) G_{14}^0 + \frac{(a_{13})^{(1)} (\hat{P}_{13})^{(1)}}{(\hat{M}_{13})^{(1)}} \left(e^{(\hat{M}_{13})^{(1)} t} - 1\right)$ | 167 |
| From which it follows that<br>$(G_{13}(t) - G_{13}^0) e^{-(\hat{M}_{13})^{(1)} t} \le \frac{(a_{13})^{(1)}}{(\hat{M}_{13})^{(1)}} \left[\left((\hat{P}_{13})^{(1)} + G_{14}^0\right) e^{\left(-\frac{(\hat{P}_{13})^{(1)} + G_{14}^0}{G_{14}^0}\right)} + (\hat{P}_{13})^{(1)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 1 | 168 |
| Analogous inequalities hold also for $G_{14}, G_{15}, T_{13}, T_{14}, T_{15}$ | |
| The operator $\mathcal{A}^{(2)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{16}(t) \le G_{16}^0 + \int_0^t \left[(a_{16})^{(2)} \left(G_{17}^0 + (\hat{P}_{16})^{(6)} e^{(\hat{M}_{16})^{(2)} s_{(16)}}\right)\right] ds_{(16)} =$<br>$\left(1 + (a_{16})^{(2)} t\right) G_{17}^0 + \frac{(a_{16})^{(2)} (\hat{P}_{16})^{(2)}}{(\hat{M}_{16})^{(2)}} \left(e^{(\hat{M}_{16})^{(2)} t} - 1\right)$ | 169 |
| From which it follows that<br>$(G_{16}(t) - G_{16}^0) e^{-(\hat{M}_{16})^{(2)} t} \le \frac{(a_{16})^{(2)}}{(\hat{M}_{16})^{(2)}} \left[\left((\hat{P}_{16})^{(2)} + G_{17}^0\right) e^{\left(-\frac{(\hat{P}_{16})^{(2)} + G_{17}^0}{G_{17}^0}\right)} + (\hat{P}_{16})^{(2)}\right]$ | 170 |
| Analogous inequalities hold also for $G_{17}, G_{18}, T_{16}, T_{17}, T_{18}$ | |
| The operator $\mathcal{A}^{(3)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{20}(t) \le G_{20}^0 + \int_0^t \left[(a_{20})^{(3)} \left(G_{21}^0 + (\hat{P}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} s_{(20)}}\right)\right] ds_{(20)} =$<br>$\left(1 + (a_{20})^{(3)} t\right) G_{21}^0 + \frac{(a_{20})^{(3)} (\hat{P}_{20})^{(3)}}{(\hat{M}_{20})^{(3)}} \left(e^{(\hat{M}_{20})^{(3)} t} - 1\right)$ | 171 |
| From which it follows that | 172 |

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| $(G_{20}(t)-G_{20}^0)e^{-(\hat{M}_{20})^{(3)}t} \le \frac{(a_{20})^{(3)}}{(\hat{M}_{20})^{(3)}}\left[((\hat{P}_{20})^{(3)}+G_{21}^0)e^{\left(-\frac{(\hat{P}_{20})^{(3)}+G_{21}^0}{G_{21}^0}\right)}+(\hat{P}_{20})^{(3)}\right]$ | |
| Analogous inequalities hold also for $G_{21}, G_{22}, T_{20}, T_{21}, T_{22}$ | |
| The operator $\mathcal{A}^{(4)}$ maps the space of functions satisfying into itself .Indeed it is obvious that $G_{24}(t) \le G_{24}^0 + \int_0^t \left[(a_{24})^{(4)}\left(G_{25}^0+(\hat{P}_{24})^{(4)}e^{(\hat{M}_{24})^{(4)}s_{(24)}}\right)\right] ds_{(24)} =$ $\left(1+(a_{24})^{(4)}t\right)G_{25}^0 + \frac{(a_{24})^{(4)}(\hat{P}_{24})^{(4)}}{(\hat{M}_{24})^{(4)}}\left(e^{(\hat{M}_{24})^{(4)}t}-1\right)$ | 173 |
| From which it follows that $(G_{24}(t)-G_{24}^0)e^{-(\hat{M}_{24})^{(4)}t} \le \frac{(a_{24})^{(4)}}{(\hat{M}_{24})^{(4)}}\left[((\hat{P}_{24})^{(4)}+G_{25}^0)e^{\left(-\frac{(\hat{P}_{24})^{(4)}+G_{25}^0}{G_{25}^0}\right)}+(\hat{P}_{24})^{(4)}\right]$ $(G_i^0)$ is as defined in the statement of theorem 4 | 174 |
| The operator $\mathcal{A}^{(5)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{28}(t) \le G_{28}^0 + \int_0^t \left[(a_{28})^{(5)}\left(G_{29}^0+(\hat{P}_{28})^{(5)}e^{(\hat{M}_{28})^{(5)}s_{(28)}}\right)\right] ds_{(28)} =$ $\left(1+(a_{28})^{(5)}t\right)G_{29}^0 + \frac{(a_{28})^{(5)}(\hat{P}_{28})^{(5)}}{(\hat{M}_{28})^{(5)}}\left(e^{(\hat{M}_{28})^{(5)}t}-1\right)$ | |
| From which it follows that $(G_{28}(t)-G_{28}^0)e^{-(\hat{M}_{28})^{(5)}t} \le \frac{(a_{28})^{(5)}}{(\hat{M}_{28})^{(5)}}\left[((\hat{P}_{28})^{(5)}+G_{29}^0)e^{\left(-\frac{(\hat{P}_{28})^{(5)}+G_{29}^0}{G_{29}^0}\right)}+(\hat{P}_{28})^{(5)}\right]$ $(G_i^0)$ is as defined in the statement of theorem 5 | 175 |
| The operator $\mathcal{A}^{(6)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{32}(t) \le G_{32}^0 + \int_0^t \left[(a_{32})^{(6)}\left(G_{33}^0+(\hat{P}_{32})^{(6)}e^{(\hat{M}_{32})^{(6)}s_{(32)}}\right)\right] ds_{(32)} =$ $\left(1+(a_{32})^{(6)}t\right)G_{33}^0 + \frac{(a_{32})^{(6)}(\hat{P}_{32})^{(6)}}{(\hat{M}_{32})^{(6)}}\left(e^{(\hat{M}_{32})^{(6)}t}-1\right)$ | 176 |
| From which it follows that $(G_{32}(t)-G_{32}^0)e^{-(\hat{M}_{32})^{(6)}t} \le \frac{(a_{32})^{(6)}}{(\hat{M}_{32})^{(6)}}\left[((\hat{P}_{32})^{(6)}+G_{33}^0)e^{\left(-\frac{(\hat{P}_{32})^{(6)}+G_{33}^0}{G_{33}^0}\right)}+(\hat{P}_{32})^{(6)}\right]$ $(G_i^0)$ is as defined in the statement of theorem 6 | 177 |

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| Analogous inequalities hold also for $G_{25}, G_{26}, T_{24}, T_{25}, T_{26}$ | |
| (d) The operator $\mathcal{A}^{(7)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{36}(t) \le G_{36}^0 + \int_0^t \left[(a_{36})^{(7)} \left(G_{37}^0 + (\hat{P}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} s_{(36)}}\right)\right] ds_{(36)} =$<br>$\left(1 + (a_{36})^{(7)} t\right) G_{37}^0 + \frac{(a_{36})^{(7)} (\hat{P}_{36})^{(7)}}{(\hat{M}_{36})^{(7)}} \left(e^{(\hat{M}_{36})^{(7)} t} - 1\right)$ | 178 |
| From which it follows that<br>$(G_{36}(t) - G_{36}^0) e^{-(\hat{M}_{36})^{(7)} t} \le \frac{(a_{36})^{(7)}}{(\hat{M}_{36})^{(7)}} \left[((\hat{P}_{36})^{(7)} + G_{37}^0) e^{\left(-\frac{(\hat{P}_{36})^{(7)} + G_{37}^0}{G_{37}^0}\right)} + (\hat{P}_{36})^{(7)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 7 | |
| The operator $\mathcal{A}^{(8)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that | |
| $G_{40}(t) \le G_{40}^0 + \int_0^t \left[(a_{40})^{(8)} \left(G_{41}^0 + (\hat{P}_{40})^{(8)} e^{(\hat{M}_{40})^{(8)} s_{(40)}}\right)\right] ds_{(40)} =$<br>$\left(1 + (a_{40})^{(8)} t\right) G_{41}^0 + \frac{(a_{40})^{(8)} (\hat{P}_{40})^{(8)}}{(\hat{M}_{40})^{(8)}} \left(e^{(\hat{M}_{40})^{(8)} t} - 1\right)$ | 180 |
| From which it follows that<br>$(G_{40}(t) - G_{40}^0) e^{-(\hat{M}_{40})^{(8)} t} \le \frac{(a_{40})^{(8)}}{(\hat{M}_{40})^{(8)}} \left[((\hat{P}_{40})^{(8)} + G_{41}^0) e^{\left(-\frac{(\hat{P}_{40})^{(8)} + G_{41}^0}{G_{41}^0}\right)} + (\hat{P}_{40})^{(8)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 8<br>Analogous inequalities hold also for $G_{41}, G_{42}, T_{40}, T_{41}, T_{42}$ | 181 |
| The operator $\mathcal{A}^{(9)}$ maps the space of functions satisfying 34,35,36 into itself .Indeed it is obvious that<br>$G_{44}(t) \le G_{44}^0 + \int_0^t \left[(a_{44})^{(9)} \left(G_{45}^0 + (\hat{P}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} s_{(44)}}\right)\right] ds_{(44)} =$<br>$\left(1 + (a_{44})^{(9)} t\right) G_{45}^0 + \frac{(a_{44})^{(9)} (\hat{P}_{44})^{(9)}}{(\hat{M}_{44})^{(9)}} \left(e^{(\hat{M}_{44})^{(9)} t} - 1\right)$ | |
| From which it follows that<br>$(G_{44}(t) - G_{44}^0) e^{-(\hat{M}_{44})^{(9)} t} \le \frac{(a_{44})^{(9)}}{(\hat{M}_{44})^{(9)}} \left[((\hat{P}_{44})^{(9)} + G_{45}^0) e^{\left(-\frac{(\hat{P}_{44})^{(9)} + G_{45}^0}{G_{45}^0}\right)} + (\hat{P}_{44})^{(9)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 9<br>Analogous inequalities hold also for $G_{45}, G_{46}, T_{44}, T_{45}, T_{46}$ | |
| It is now sufficient to take $\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}} < 1$ and to choose | 182 |

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| $(\hat{P}_{13})^{(1)}$ and $(\hat{Q}_{13})^{(1)}$ large to have | |
| $\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}\left[(\hat{P}_{13})^{(1)}+\left((\hat{P}_{13})^{(1)}+G_j^0\right)e^{-\left(\frac{(\hat{P}_{13})^{(1)}+G_j^0}{G_j^0}\right)}\right]\leq(\hat{P}_{13})^{(1)}$ | 183 |
| $\frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}}\left[\left((\hat{Q}_{13})^{(1)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{13})^{(1)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{13})^{(1)}\right]\leq(\hat{Q}_{13})^{(1)}$ | 184 |
| In order that the operator $\mathcal{A}^{(1)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(1)}$ is a contraction with respect to the metric<br>$d\left(\left(G^{(1)},T^{(1)}\right),\left(G^{(2)},T^{(2)}\right)\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\hat{M}_{13})^{(1)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\hat{M}_{13})^{(1)}t}\}$ | 185 |
| Indeed if we denote<br>**Definition of** $\tilde{G},\tilde{T}:\left(\tilde{G},\tilde{T}\right)=\mathcal{A}^{(1)}(G,T)$<br>It results<br>$\left\|\tilde{G}_{13}^{(1)}-\tilde{G}_i^{(2)}\right\|\leq\int_0^t(a_{13})^{(1)}\left\|G_{14}^{(1)}-G_{14}^{(2)}\right\|e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{(\hat{M}_{13})^{(1)}s_{(13)}}ds_{(13)}+$<br>$\int_0^t\{(a_{13}')^{(1)}\left\|G_{13}^{(1)}-G_{13}^{(2)}\right\|e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{-(\hat{M}_{13})^{(1)}s_{(13)}}+$<br>$(a_{13}'')^{(1)}\left(T_{14}^{(1)},s_{(13)}\right)\left\|G_{13}^{(1)}-G_{13}^{(2)}\right\|e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{(\hat{M}_{13})^{(1)}s_{(13)}}+$<br>$G_{13}^{(2)}\left\\|(a_{13}'')^{(1)}\left(T_{14}^{(1)},s_{(13)}\right)-(a_{13}'')^{(1)}\left(T_{14}^{(2)},s_{(13)}\right)\right\|\ e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{(\hat{M}_{13})^{(1)}s_{(13)}}\}ds_{(13)}$<br>Where $s_{(13)}$ represents integrand that is integrated over the interval $[0,t]$<br>From the hypotheses it follows | |
| $\left\|G^{(1)}-G^{(2)}\right\|e^{-(\hat{M}_{13})^{(1)}t}\leq$<br>$\frac{1}{(\hat{M}_{13})^{(1)}}\left((a_{13})^{(1)}+(a_{13}')^{(1)}+(\hat{A}_{13})^{(1)}+(\hat{P}_{13})^{(1)}(\hat{k}_{13})^{(1)}\right)d\left(\left(G^{(1)},T^{(1)};\ G^{(2)},T^{(2)}\right)\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 186 |
| **Remark 1:** The fact that we supposed $(a_{13}'')^{(1)}$ and $(b_{13}'')^{(1)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by | |

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| $(\widehat{P}_{13})^{(1)}e^{(\widehat{M}_{13})^{(1)}t}$ *and* $(\widehat{Q}_{13})^{(1)}e^{(\widehat{M}_{13})^{(1)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(1)}$ and $(b_i^{''})^{(1)}, i = 13,14,15$ depend only on $T_{14}$ and respectively on $G(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 2:** There does not exist any $t$ where $G_i(t) = 0\ and\ T_i(t) = 0$<br><br>From 19 to 24 it results<br><br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(1)} - (a_i^{''})^{(1)}(T_{14}(s_{(13)}), s_{(13)})\} ds_{(13)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(1)}t\right)} > 0$ for t > 0 | |
| **Definition of** $\left((\widehat{M}_{13})^{(1)}\right)_1, \left((\widehat{M}_{13})^{(1)}\right)_2\ and\ \left((\widehat{M}_{13})^{(1)}\right)_3$ :<br><br>**Remark 3**: if $G_{13}$ is bounded, the same property have also $G_{14}\ and\ G_{15}$ . indeed if<br><br>$G_{13} < (\widehat{M}_{13})^{(1)}$ it follows $\frac{dG_{14}}{dt} \leq \left((\widehat{M}_{13})^{(1)}\right)_1 - (a_{14}^{'})^{(1)}G_{14}$ and by integrating<br><br>$G_{14} \leq \left((\widehat{M}_{13})^{(1)}\right)_2 = G_{14}^0 + 2(a_{14})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_1/(a_{14}^{'})^{(1)}$<br><br>In the same way , one can obtain<br><br>$G_{15} \leq \left((\widehat{M}_{13})^{(1)}\right)_3 = G_{15}^0 + 2(a_{15})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_2/(a_{15}^{'})^{(1)}$<br><br>If $G_{14}\ or\ G_{15}$ is bounded, the same property follows for $G_{13}$ , $G_{15}$ and $G_{13}$ , $G_{14}$ respectively. | 187 |
| **Remark 4:** If $G_{13}$ *is* bounded, from below, the same property holds for $G_{14}\ and\ G_{15}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{14}$ is bounded from below. | 188 |
| **Remark 5:** If $T_{13}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(1)}(G(t), t)) = (b_{14}^{'})^{(1)}$ then $T_{14} \to \infty$.<br><br>**Definition of** $(m)^{(1)}$ and $\varepsilon_1$ :<br><br>Indeed let $t_1$ be so that for $t > t_1$<br><br>$(b_{14})^{(1)} - (b_i^{''})^{(1)}(G(t), t) < \varepsilon_1, T_{13}(t) > (m)^{(1)}$ | 189 |
| Then $\frac{dT_{14}}{dt} \geq (a_{14})^{(1)}(m)^{(1)} - \varepsilon_1 T_{14}$ which leads to<br><br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{\varepsilon_1}\right)(1 - e^{-\varepsilon_1 t}) + T_{14}^0 e^{-\varepsilon_1 t}$ If we take t such that $e^{-\varepsilon_1 t} = \frac{1}{2}$ it results<br><br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{2}\right),\quad t = log\frac{2}{\varepsilon_1}$ By taking now $\varepsilon_1$ sufficiently small one sees that $T_{14}$ is unbounded.<br>The same property holds for $T_{15}$ if $\lim_{t\to\infty}(b_{15}^{''})^{(1)}(G(t), t) = (b_{15}^{'})^{(1)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |

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| It is now sufficient to take $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}$ , $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ and to choose<br>$(\hat{P}_{16})^{(2)}$ *and* $(\hat{Q}_{16})^{(2)}$ large to have | 190 |
| $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[(\hat{P}_{16})^{(2)} + \left((\hat{P}_{16})^{(2)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{16})^{(2)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{16})^{(2)}$ | 191 |
| $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[\left((\hat{Q}_{16})^{(2)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{16})^{(2)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{16})^{(2)}\right] \leq (\hat{Q}_{16})^{(2)}$ | 192 |
| In order that the operator $\mathcal{A}^{(2)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 193 |
| The operator $\mathcal{A}^{(2)}$ is a contraction with respect to the metric<br>$d\left(((G_{19})^{(1)}, (T_{19})^{(1)}), ((G_{19})^{(2)}, (T_{19})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\hat{M}_{16})^{(2)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\hat{M}_{16})^{(2)}t}\}$ | 194 |
| Indeed if we denote<br>**Definition of** $\widetilde{G_{19}}, \widetilde{T_{19}}$ : $\left(\widetilde{G_{19}}, \widetilde{T_{19}}\right) = \mathcal{A}^{(2)}(G_{19}, T_{19})$ | 195 |
| It results<br>$\left\|\tilde{G}_{16}^{(1)} - \tilde{G}_i^{(2)}\right\| \leq \int_0^t (a_{16})^{(2)}\left\|G_{17}^{(1)} - G_{17}^{(2)}\right\|e^{-(\hat{M}_{16})^{(2)}s_{(16)}}e^{(\hat{M}_{16})^{(2)}s_{(16)}}\,ds_{(16)} +$<br>$\int_0^t\{(a_{16}')^{(2)}\left\|G_{16}^{(1)} - G_{16}^{(2)}\right\|e^{-(\hat{M}_{16})^{(2)}s_{(16)}}e^{-(\hat{M}_{16})^{(2)}s_{(16)}} +$<br>$(a_{16}'')^{(2)}\left(T_{17}^{(1)}, s_{(16)}\right)\left\|G_{16}^{(1)} - G_{16}^{(2)}\right\|e^{-(\hat{M}_{16})^{(2)}s_{(16)}}e^{(\hat{M}_{16})^{(2)}s_{(16)}} +$<br>$G_{16}^{(2)}\left\|(a_{16}'')^{(2)}\left(T_{17}^{(1)}, s_{(16)}\right) - (a_{16}'')^{(2)}\left(T_{17}^{(2)}, s_{(16)}\right)\right\| \; e^{-(\hat{M}_{16})^{(2)}s_{(16)}}e^{(\hat{M}_{16})^{(2)}s_{(16)}}\}ds_{(16)}$ | 196 |
| Where $s_{(16)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | 197 |
| $\left\|(G_{19})^{(1)} - (G_{19})^{(2)}\right\|e^{-(\hat{M}_{16})^{(2)}t} \leq$<br>$\frac{1}{(\hat{M}_{16})^{(2)}}\left((a_{16})^{(2)} + (a_{16}')^{(2)} + (\hat{A}_{16})^{(2)} + (\hat{P}_{16})^{(2)}(\hat{k}_{16})^{(2)}\right)d\left(((G_{19})^{(1)}, (T_{19})^{(1)}; (G_{19})^{(2)}, (T_{19})^{(2)})\right)$ | |
| And analogous inequalities for $G_i$ and $T_i$. Taking into account the hypothesis the result follows | 198 |
| **Remark 6:** The fact that we supposed $(a_{16}'')^{(2)}$ and $(b_{16}'')^{(2)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate | 199 |

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| condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{16})^{(2)}e^{(\widehat{M}_{16})^{(2)}t}$ and $(\widehat{Q}_{16})^{(2)}e^{(\widehat{M}_{16})^{(2)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}, i = 16,17,18$ depend only on $T_{17}$ and respectively on $(G_{19})$(and not on t) and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 7:** There does not exist any t where $G_i(t) = 0$ and $T_i(t) = 0$<br><br>it results<br><br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(2)} - (a_i^{''})^{(2)}(T_{17}(s_{(16)}), s_{(16)})\} ds_{(16)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{(-(b_i^{'})^{(2)}t)} > 0$ for $t > 0$ | 200 |
| **Definition of** $\left((\widehat{M}_{16})^{(2)}\right)_1, \left((\widehat{M}_{16})^{(2)}\right)_2$ and $\left((\widehat{M}_{16})^{(2)}\right)_3$ :<br><br>**Remark 8:** if $G_{16}$ is bounded, the same property have also $G_{17}$ and $G_{18}$ . indeed if<br><br>$G_{16} < (\widehat{M}_{16})^{(2)}$ it follows $\frac{dG_{17}}{dt} \leq \left((\widehat{M}_{16})^{(2)}\right)_1 - (a_{17}^{'})^{(2)}G_{17}$ and by integrating<br><br>$G_{17} \leq \left((\widehat{M}_{16})^{(2)}\right)_2 = G_{17}^0 + 2(a_{17})^{(2)}\left((\widehat{M}_{16})^{(2)}\right)_1/(a_{17}^{'})^{(2)}$<br><br>In the same way , one can obtain<br><br>$G_{18} \leq \left((\widehat{M}_{16})^{(2)}\right)_3 = G_{18}^0 + 2(a_{18})^{(2)}\left((\widehat{M}_{16})^{(2)}\right)_2/(a_{18}^{'})^{(2)}$<br><br>If $G_{17}$ or $G_{18}$ is bounded, the same property follows for $G_{16}$ , $G_{18}$ and $G_{16}$ , $G_{17}$ respectively. | 201 |
| **Remark 9**: If $G_{16}$ is bounded, from below, the same property holds for $G_{17}$ and $G_{18}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{17}$ is bounded from below. | 202 |
| **Remark 10:** If $T_{16}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(2)}((G_{19})(t), t)) = (b_{17}^{'})^{(2)}$ then $T_{17} \to \infty$.<br><br>**Definition of** $(m)^{(2)}$ and $\varepsilon_2$ :<br><br>Indeed let $t_2$ be so that for $t > t_2$<br><br>$(b_{17})^{(2)} - (b_i^{''})^{(2)}((G_{19})(t), t) < \varepsilon_2, T_{16}(t) > (m)^{(2)}$ | 203 |
| Then $\frac{dT_{17}}{dt} \geq (a_{17})^{(2)}(m)^{(2)} - \varepsilon_2 T_{17}$ which leads to<br><br>$T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{\varepsilon_2}\right)(1 - e^{-\varepsilon_2 t}) + T_{17}^0 e^{-\varepsilon_2 t}$ If we take t such that $e^{-\varepsilon_2 t} = \frac{1}{2}$ it results | 204 |
| $T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{2}\right)$, $t = \log\frac{2}{\varepsilon_2}$ By taking now $\varepsilon_2$ sufficiently small one sees that $T_{17}$ is unbounded.<br>The same property holds for $T_{18}$ if $\lim_{t\to\infty}(b_{18}^{''})^{(2)}\left((G_{19})(t), t\right) = (b_{18}^{'})^{(2)}$ | 205 |

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| We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |
| It is now sufficient to take $\frac{(a_i)^{(3)}}{(\widehat{M}_{20})^{(3)}}$ , $\frac{(b_i)^{(3)}}{(\widehat{M}_{20})^{(3)}} < 1$ and to choose<br>$(\widehat{P}_{20})^{(3)}$ and $(\widehat{Q}_{20})^{(3)}$ large to have | 207 |
| $\frac{(a_i)^{(3)}}{(\widehat{M}_{20})^{(3)}}\left[(\widehat{P}_{20})^{(3)} + \left((\hat{P}_{20})^{(3)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{20})^{(3)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{20})^{(3)}$ | 208 |
| $\frac{(b_i)^{(3)}}{(\widehat{M}_{20})^{(3)}}\left[\left((\hat{Q}_{20})^{(3)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{20})^{(3)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{20})^{(3)}\right] \leq (\hat{Q}_{20})^{(3)}$ | 209 |
| In order that the operator $\mathcal{A}^{(3)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 210 |
| The operator $\mathcal{A}^{(3)}$ is a contraction with respect to the metric<br>$d\left(((G_{23})^{(1)}, (T_{23})^{(1)}), ((G_{23})^{(2)}, (T_{23})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{20})^{(3)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{20})^{(3)}t}\}$ | 211 |
| Indeed if we denote<br>**Definition of** $\widetilde{G_{23}}, \widetilde{T_{23}} :(\widetilde{(G_{23})}, \widetilde{(T_{23})}) = \mathcal{A}^{(3)}((G_{23}), (T_{23}))$ | 212 |
| It results<br>$\left\|\tilde{G}_{20}^{(1)} - \tilde{G}_i^{(2)}\right\| \leq \int_0^t (a_{20})^{(3)}\left\|G_{21}^{(1)} - G_{21}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{(\widehat{M}_{20})^{(3)}s_{(20)}}\,ds_{(20)} +$<br>$\int_0^t\{(a'_{20})^{(3)}\left\|G_{20}^{(1)} - G_{20}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{-(\widehat{M}_{20})^{(3)}s_{(20)}} +$<br>$(a''_{20})^{(3)}\left(T_{21}^{(1)}, s_{(20)}\right)\left\|G_{20}^{(1)} - G_{20}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{(\widehat{M}_{20})^{(3)}s_{(20)}} +$<br>$G_{20}^{(2)}\left\|(a''_{20})^{(3)}\left(T_{21}^{(1)}, s_{(20)}\right) - (a''_{20})^{(3)}\left(T_{21}^{(2)}, s_{(20)}\right)\right\| \; e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{(\widehat{M}_{20})^{(3)}s_{(20)}}\}ds_{(20)}$<br>Where $s_{(20)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | 213 |
| $\left\|G_{23}{}^{(1)} - G_{23}{}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}t} \leq$<br>$\frac{1}{(\widehat{M}_{20})^{(3)}}\left((a_{20})^{(3)} + (a'_{20})^{(3)} + (\widehat{A}_{20})^{(3)} + (\widehat{P}_{20})^{(3)}(\widehat{k}_{20})^{(3)}\right)d\left(((G_{23})^{(1)}, (T_{23})^{(1)};\ (G_{23})^{(2)}, (T_{23})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 214 |

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| **Remark 11:** The fact that we supposed $(a_{20}^{''})^{(3)}$ and $(b_{20}^{''})^{(3)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}$ $and$ $(\widehat{Q}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(3)}$ and $(b_i^{''})^{(3)}, i = 20,21,22$ depend only on $T_{21}$ and respectively on $(G_{23})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 215 |
| **Remark 12:** There does not exist any $t$ where $G_i(t) = 0\ and\ T_i(t) = 0$<br><br>it results<br><br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(3)} - (a_i^{''})^{(3)}(T_{21}(s_{(20)}), s_{(20)})\} ds_{(20)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{(-(b_i^{'})^{(3)}t)} > 0$ for t > 0 | 216 |
| **Definition of** $\left((\widehat{M}_{20})^{(3)}\right)_1, \left((\widehat{M}_{20})^{(3)}\right)_2\ and\ \left((\widehat{M}_{20})^{(3)}\right)_3$ :<br><br>**Remark 13:** if $G_{20}$ is bounded, the same property have also $G_{21}\ and\ G_{22}$ . indeed if<br><br>$G_{20} < (\widehat{M}_{20})^{(3)}$ it follows $\frac{dG_{21}}{dt} \leq \left((\widehat{M}_{20})^{(3)}\right)_1 - (a_{21}^{'})^{(3)}G_{21}$ and by integrating<br><br>$G_{21} \leq \left((\widehat{M}_{20})^{(3)}\right)_2 = G_{21}^0 + 2(a_{21})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_1/(a_{21}^{'})^{(3)}$<br><br>In the same way , one can obtain<br><br>$G_{22} \leq \left((\widehat{M}_{20})^{(3)}\right)_3 = G_{22}^0 + 2(a_{22})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_2/(a_{22}^{'})^{(3)}$<br><br>If $G_{21}\ or\ G_{22}$ is bounded, the same property follows for $G_{20}$ , $G_{22}$ and $G_{20}$ , $G_{21}$ respectively. | 217 |
| **Remark 14:** If $G_{20}$ $is$ bounded, from below, the same property holds for $G_{21}\ and\ G_{22}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{21}$ is bounded from below. | 218 |
| **Remark 15:** If $T_{20}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(3)}\left((G_{23})(t), t\right)) = (b_{21}^{'})^{(3)}$ then $T_{21} \to \infty$.<br><br>**Definition of** $(m)^{(3)}$ and $\varepsilon_3$ :<br><br>Indeed let $t_3$ be so that for $t > t_3$<br><br>$(b_{21})^{(3)} - (b_i^{''})^{(3)}\left((G_{23})(t), t\right) < \varepsilon_3, T_{20}(t) > (m)^{(3)}$ | 219 |
| Then $\frac{dT_{21}}{dt} \geq (a_{21})^{(3)}(m)^{(3)} - \varepsilon_3 T_{21}$ which leads to<br><br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{\varepsilon_3}\right)(1 - e^{-\varepsilon_3 t}) + T_{21}^0 e^{-\varepsilon_3 t}$ If we take t such that $e^{-\varepsilon_3 t} = \frac{1}{2}$ it results<br><br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{2}\right)$, $t = log\frac{2}{\varepsilon_3}$ By taking now $\varepsilon_3$ sufficiently small one sees that $T_{21}$ is unbounded. | 220 |

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| The same property holds for $T_{22}$ if $\lim_{t\to\infty}(b_{22}'')^{(3)}\left((G_{23})(t),t\right)=(b_{22}')^{(3)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |
| It is now sufficient to take $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}$ , $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}<1$ and to choose<br>$(\widehat{P}_{24})^{(4)}$ and $(\widehat{Q}_{24})^{(4)}$ large to have | 221 |
| $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[(\widehat{P}_{24})^{(4)}+\left((\hat{P}_{24})^{(4)}+G_j^0\right)e^{-\left(\frac{(\hat{P}_{24})^{(4)}+G_j^0}{G_j^0}\right)}\right]\le(\hat{P}_{24})^{(4)}$ | 222 |
| $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[\left((\hat{Q}_{24})^{(4)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{24})^{(4)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{24})^{(4)}\right]\le(\hat{Q}_{24})^{(4)}$ | 223 |
| In order that the operator $\mathcal{A}^{(4)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 224 |
| The operator $\mathcal{A}^{(4)}$ is a contraction with respect to the metric<br>$d\left(((G_{27})^{(1)},(T_{27})^{(1)}),((G_{27})^{(2)},(T_{27})^{(2)})\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{24})^{(4)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{24})^{(4)}t}\}$<br>Indeed if we denote<br>**Definition of** $(\widetilde{G_{27}}),(\widetilde{T_{27}})$ : $\left((\widetilde{G_{27}}),(\widetilde{T_{27}})\right)=\mathcal{A}^{(4)}((G_{27}),(T_{27}))$<br>It results<br>$\left\|\tilde{G}_{24}^{(1)}-\tilde{G}_i^{(2)}\right\|\le\int_0^t(a_{24})^{(4)}\left\|G_{25}^{(1)}-G_{25}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}ds_{(24)}+$<br>$\int_0^t\{(a_{24}')^{(4)}\left\|G_{24}^{(1)}-G_{24}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}+$<br>$(a_{24}'')^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)\left\|G_{24}^{(1)}-G_{24}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}+$<br>$G_{24}^{(2)}\left\\|(a_{24}'')^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)-(a_{24}'')^{(4)}\left(T_{25}^{(2)},s_{(24)}\right)\right\\|\ e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}\}ds_{(24)}$<br>Where $s_{(24)}$ represents integrand that is integrated over the interval $[0,t]$<br>From the hypotheses on Equations it follows | 225 |
| $\left\|(G_{27})^{(1)}-(G_{27})^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}t}\le$<br>$\frac{1}{(\widehat{M}_{24})^{(4)}}\left((a_{24})^{(4)}+(a_{24}')^{(4)}+(\widehat{A}_{24})^{(4)}+(\widehat{P}_{24})^{(4)}(\widehat{k}_{24})^{(4)}\right)d\left(((G_{27})^{(1)},(T_{27})^{(1)};\ (G_{27})^{(2)},(T_{27})^{(2)})\right)$ | 226 |

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| And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | |
| **Remark 16:** The fact that we supposed $(a_{24}^{''})^{(4)}$ and $(b_{24}^{''})^{(4)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}$ *and* $(\widehat{Q}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(4)}$ and $(b_i^{''})^{(4)}, i = 24,25,26$ depend only on $T_{25}$ and respectively on $(G_{27})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 227 |
| **Remark 17:** There does not exist any $t$ where $G_i\,(t) = 0\ and\ T_i\,(t) = 0$<br>it results<br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t\{(a_i^{'})^{(4)}-(a_i^{''})^{(4)}(T_{25}(s_{(24)}),s_{(24)})\}ds_{(24)}\right]} \geq 0$<br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(4)}t\right)} > 0$ for t > 0 | 228 |
| **Definition of** $\left((\widehat{M}_{24})^{(4)}\right)_1, \left((\widehat{M}_{24})^{(4)}\right)_2\ and\ \left((\widehat{M}_{24})^{(4)}\right)_3$ :<br>**Remark 18:** if $G_{24}$ is bounded, the same property have also $G_{25}\ and\ G_{26}$ . indeed if<br>$G_{24} < (\widehat{M}_{24})^{(4)}$ it follows $\frac{dG_{25}}{dt} \leq \left((\widehat{M}_{24})^{(4)}\right)_1 - (a_{25}^{'})^{(4)}G_{25}$ and by integrating<br>$G_{25} \leq \left((\widehat{M}_{24})^{(4)}\right)_2 = G_{25}^0 + 2(a_{25})^{(4)}\left((\widehat{M}_{24})^{(4)}\right)_1/(a_{25}^{'})^{(4)}$<br>In the same way , one can obtain<br>$G_{26} \leq \left((\widehat{M}_{24})^{(4)}\right)_3 = G_{26}^0 + 2(a_{26})^{(4)}\left((\widehat{M}_{24})^{(4)}\right)_2/(a_{26}^{'})^{(4)}$<br>If $G_{25}\ or\ G_{26}$ is bounded, the same property follows for $G_{24}$ , $G_{26}$ and $G_{24}$ , $G_{25}$ respectively. | 229 |
| **Remark 19:** If $G_{24}$ *is* bounded, from below, the same property holds for $G_{25}\ and\ G_{26}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{25}$ is bounded from below. | 230 |
| **Remark 20:** If $T_{24}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(4)}((G_{27})(t),t)) = (b_{25}^{'})^{(4)}$ then $T_{25} \to \infty$.<br>**Definition of** $(m)^{(4)}$ and $\varepsilon_4$ :<br>Indeed let $t_4$ be so that for $t > t_4$<br>$(b_{25})^{(4)} - (b_i^{''})^{(4)}((G_{27})(t),t) < \varepsilon_4, T_{24}\,(t) > (m)^{(4)}$ | 231 |
| Then $\frac{dT_{25}}{dt} \geq (a_{25})^{(4)}(m)^{(4)} - \varepsilon_4 T_{25}$ which leads to<br>$T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{\varepsilon_4}\right)(1 - e^{-\varepsilon_4 t}) + T_{25}^0 e^{-\varepsilon_4 t}$ If we take t such that $e^{-\varepsilon_4 t} = \frac{1}{2}$ it results | 232 |

$T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{2}\right), \quad t = log\frac{2}{\varepsilon_4}$ By taking now $\varepsilon_4$ sufficiently small one sees that $T_{25}$ is unbounded. The same property holds for $T_{26}$ if $\lim_{t\to\infty}(b_{26}^{''})^{(4)}\left((G_{27})(t),t\right) = (b_{26}^{'})^{(4)}$

We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 42

Analogous inequalities hold also for $G_{29}, G_{30}, T_{28}, T_{29}, T_{30}$

It is now sufficient to take $\frac{(a_i)^{(5)}}{(\hat{M}_{28})^{(5)}}, \frac{(b_i)^{(5)}}{(\hat{M}_{28})^{(5)}} < 1$ and to choose

$(\hat{P}_{28})^{(5)}$ and $(\hat{Q}_{28})^{(5)}$ large to have

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$$\frac{(a_i)^{(5)}}{(\hat{M}_{28})^{(5)}}\left[(\hat{P}_{28})^{(5)} + \left((\hat{P}_{28})^{(5)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{28})^{(5)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{28})^{(5)} \quad (234)$$

$$\frac{(b_i)^{(5)}}{(\hat{M}_{28})^{(5)}}\left[\left((\hat{Q}_{28})^{(5)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{28})^{(5)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{28})^{(5)}\right] \leq (\hat{Q}_{28})^{(5)} \quad (235)$$

In order that the operator $\mathcal{A}^{(5)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself

The operator $\mathcal{A}^{(5)}$ is a contraction with respect to the metric (236)

$$d\left(((G_{31})^{(1)},(T_{31})^{(1)}),((G_{31})^{(2)},(T_{31})^{(2)})\right) =$$

$$\sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t) - G_i^{(2)}(t)\right|e^{-(\hat{M}_{28})^{(5)}t}, \max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t) - T_i^{(2)}(t)\right|e^{-(\hat{M}_{28})^{(5)}t}\}$$

Indeed if we denote

**Definition of** $(\widetilde{G_{31}}),(\widetilde{T_{31}})$ : $\left((\widetilde{G_{31}}),(\widetilde{T_{31}})\right) = \mathcal{A}^{(5)}\left((G_{31}),(T_{31})\right)$

It results

$$\left|\tilde{G}_{28}^{(1)} - \tilde{G}_i^{(2)}\right| \leq \int_0^t (a_{28})^{(5)}\left|G_{29}^{(1)} - G_{29}^{(2)}\right|e^{-(\hat{M}_{28})^{(5)}s_{(28)}}e^{(\hat{M}_{28})^{(5)}s_{(28)}}\,ds_{(28)} +$$

$$\int_0^t\{(a_{28}^{'})^{(5)}\left|G_{28}^{(1)} - G_{28}^{(2)}\right|e^{-(\hat{M}_{28})^{(5)}s_{(28)}}e^{-(\hat{M}_{28})^{(5)}s_{(28)}} +$$

$$(a_{28}^{''})^{(5)}\left(T_{29}^{(1)},s_{(28)}\right)\left|G_{28}^{(1)} - G_{28}^{(2)}\right|e^{-(\hat{M}_{28})^{(5)}s_{(28)}}e^{(\hat{M}_{28})^{(5)}s_{(28)}} +$$

$$G_{28}^{(2)}\left|(a_{28}^{''})^{(5)}\left(T_{29}^{(1)},s_{(28)}\right) - (a_{28}^{''})^{(5)}\left(T_{29}^{(2)},s_{(28)}\right)\right|\ e^{-(\hat{M}_{28})^{(5)}s_{(28)}}e^{(\hat{M}_{28})^{(5)}s_{(28)}}\}ds_{(28)}$$

Where $s_{(28)}$ represents integrand that is integrated over the interval [0, t]

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| From the hypotheses on it follows | |
| $\left\|(G_{31})^{(1)} - (G_{31})^{(2)}\right\| e^{-(\widehat{M}_{28})^{(5)}t} \leq$ $\frac{1}{(\widehat{M}_{28})^{(5)}}\left((a_{28})^{(5)} + (a'_{28})^{(5)} + (\widehat{A}_{28})^{(5)} + (\widehat{P}_{28})^{(5)}(\widehat{k}_{28})^{(5)}\right) d\left(((G_{31})^{(1)}, (T_{31})^{(1)};\ (G_{31})^{(2)}, (T_{31})^{(2)})\right)$ And analogous inequalities for $G_i$ and $T_i$. Taking into account the hypothesis the result follows | 237 |
| **Remark 21:** The fact that we supposed $(a''_{28})^{(5)}$ and $(b''_{28})^{(5)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{28})^{(5)} e^{(\widehat{M}_{28})^{(5)}t}$ and $(\widehat{Q}_{28})^{(5)} e^{(\widehat{M}_{28})^{(5)}t}$ respectively of $\mathbb{R}_+$. If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a''_i)^{(5)}$ and $(b''_i)^{(5)}, i = 28,29,30$ depend only on $T_{29}$ and respectively on $(G_{31})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 238 |
| **Remark 22:** There does not exist any $t$ where $G_i(t) = 0$ and $T_i(t) = 0$ it results $G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a'_i)^{(5)} - (a''_i)^{(5)}(T_{29}(s_{(28)}), s_{(28)})\} ds_{(28)}\right]} \geq 0$ $T_i(t) \geq T_i^0 e^{\left(-(b'_i)^{(5)}t\right)} > 0$ for t > 0 | 239 |
| **Definition of** $\left((\widehat{M}_{28})^{(5)}\right)_1, \left((\widehat{M}_{28})^{(5)}\right)_2$ and $\left((\widehat{M}_{28})^{(5)}\right)_3$ : **Remark 23:** if $G_{28}$ is bounded, the same property have also $G_{29}$ and $G_{30}$ . indeed if $G_{28} < (\widehat{M}_{28})^{(5)}$ it follows $\frac{dG_{29}}{dt} \leq \left((\widehat{M}_{28})^{(5)}\right)_1 - (a'_{29})^{(5)} G_{29}$ and by integrating $G_{29} \leq \left((\widehat{M}_{28})^{(5)}\right)_2 = G_{29}^0 + 2(a_{29})^{(5)}\left((\widehat{M}_{28})^{(5)}\right)_1 / (a'_{29})^{(5)}$ In the same way , one can obtain $G_{30} \leq \left((\widehat{M}_{28})^{(5)}\right)_3 = G_{30}^0 + 2(a_{30})^{(5)}\left((\widehat{M}_{28})^{(5)}\right)_2 / (a'_{30})^{(5)}$ If $G_{29}$ or $G_{30}$ is bounded, the same property follows for $G_{28}$ , $G_{30}$ and $G_{28}$ , $G_{29}$ respectively. | 240 |
| **Remark 24:** If $G_{28}$ is bounded, from below, the same property holds for $G_{29}$ and $G_{30}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{29}$ is bounded from below. | 241 |

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| **Remark 25:** If $T_{28}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(5)}((G_{31})(t),t)) = (b_{29}^{'})^{(5)}$ then $T_{29} \to \infty$.<br><br>**Definition of** $(m)^{(5)}$ and $\varepsilon_5$ :<br><br>Indeed let $t_5$ be so that for $t > t_5$<br><br>$(b_{29})^{(5)} - (b_i^{''})^{(5)}((G_{31})(t),t) < \varepsilon_5, T_{28}(t) > (m)^{(5)}$ | 242 |
| Then $\frac{dT_{29}}{dt} \geq (a_{29})^{(5)}(m)^{(5)} - \varepsilon_5 T_{29}$ which leads to<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{\varepsilon_5}\right)(1-e^{-\varepsilon_5 t}) + T_{29}^0 e^{-\varepsilon_5 t}$ If we take t such that $e^{-\varepsilon_5 t} = \frac{1}{2}$ it results<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{2}\right), \quad t = log\frac{2}{\varepsilon_5}$ By taking now $\varepsilon_5$ sufficiently small one sees that $T_{29}$ is unbounded.<br>The same property holds for $T_{30}$ if $\lim_{t\to\infty}(b_{30}^{''})^{(5)}\left((G_{31})(t),t\right) = (b_{30}^{'})^{(5)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations<br><br>Analogous inequalities hold also for $G_{33}, G_{34}, T_{32}, T_{33}, T_{34}$ | 243 |
| It is now sufficient to take $\frac{(a_i)^{(6)}}{(\hat{M}_{32})^{(6)}}, \frac{(b_i)^{(6)}}{(\hat{M}_{32})^{(6)}} < 1$ and to choose<br><br>$(\hat{P}_{32})^{(6)}$ and $(\hat{Q}_{32})^{(6)}$ large to have | 244 |
| $\frac{(a_i)^{(6)}}{(\hat{M}_{32})^{(6)}}\left[(\hat{P}_{32})^{(6)} + \left((\hat{P}_{32})^{(6)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{32})^{(6)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{32})^{(6)}$ | 245 |
| $\frac{(b_i)^{(6)}}{(\hat{M}_{32})^{(6)}}\left[\left((\hat{Q}_{32})^{(6)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{32})^{(6)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{32})^{(6)}\right] \leq (\hat{Q}_{32})^{(6)}$ | 246 |
| In order that the operator $\mathcal{A}^{(6)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(6)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{35})^{(1)},(T_{35})^{(1)}),((G_{35})^{(2)},(T_{35})^{(2)})\right) =$<br><br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t) - G_i^{(2)}(t)\right|e^{-(\hat{M}_{32})^{(6)}t}, \max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t) - T_i^{(2)}(t)\right|e^{-(\hat{M}_{32})^{(6)}t}\}$<br><br>Indeed if we denote<br><br>**Definition of** $\widetilde{(G_{35})}, \widetilde{(T_{35})}$ : $\left(\widetilde{(G_{35})}, \widetilde{(T_{35})}\right) = \mathcal{A}^{(6)}\left((G_{35}),(T_{35})\right)$<br><br>It results | 247 |

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| $\left\|\tilde{G}_{32}^{(1)}-\tilde{G}_{i}^{(2)}\right\| \leq \int_0^t (a_{32})^{(6)} \left\|G_{33}^{(1)}-G_{33}^{(2)}\right\| e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} e^{(\widehat{M}_{32})^{(6)}s_{(32)}}\, ds_{(32)} +$<br>$\int_0^t \{(a_{32}')^{(6)} \left\|G_{32}^{(1)}-G_{32}^{(2)}\right\| e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} +$<br>$(a_{32}'')^{(6)}\left(T_{33}^{(1)}, s_{(32)}\right)\left\|G_{32}^{(1)}-G_{32}^{(2)}\right\| e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} e^{(\widehat{M}_{32})^{(6)}s_{(32)}} +$<br>$G_{32}^{(2)} \left\|(a_{32}'')^{(6)}\left(T_{33}^{(1)}, s_{(32)}\right) - (a_{32}'')^{(6)}\left(T_{33}^{(2)}, s_{(32)}\right)\right\| \; e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} e^{(\widehat{M}_{32})^{(6)}s_{(32)}}\} ds_{(32)}$<br>Where $s_{(32)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | |
| $\left\|(G_{35})^{(1)}-(G_{35})^{(2)}\right\| e^{-(\widehat{M}_{32})^{(6)}t} \leq$<br>$\frac{1}{(\widehat{M}_{32})^{(6)}}\left((a_{32})^{(6)} + (a_{32}')^{(6)} + (\widehat{A}_{32})^{(6)} + (\widehat{P}_{32})^{(6)}(\widehat{k}_{32})^{(6)}\right) d\left(((G_{35})^{(1)}, (T_{35})^{(1)};\ (G_{35})^{(2)}, (T_{35})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 248 |
| **Remark 26:** The fact that we supposed $(a_{32}'')^{(6)}$ and $(b_{32}'')^{(6)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{32})^{(6)} e^{(\widehat{M}_{32})^{(6)}t}$ *and* $(\widehat{Q}_{32})^{(6)} e^{(\widehat{M}_{32})^{(6)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i'')^{(6)}$ and $(b_i'')^{(6)}, i = 32,33,34$ depend only on $T_{33}$ and respectively on $(G_{35})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 249 |
| **Remark 27:** There does not exist any $t$ where $G_i(t) = 0$ *and* $T_i(t) = 0$<br>it results<br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i')^{(6)} - (a_i'')^{(6)}(T_{33}(s_{(32)}), s_{(32)})\} ds_{(32)}\right]} \geq 0$<br>$T_i(t) \geq T_i^0 e^{\left(-(b_i')^{(6)}t\right)} > 0$ for t > 0 | 250 |
| **Definition of** $\left((\widehat{M}_{32})^{(6)}\right)_1, \left((\widehat{M}_{32})^{(6)}\right)_2$ *and* $\left((\widehat{M}_{32})^{(6)}\right)_3$ :<br>**Remark 28:** if $G_{32}$ is bounded, the same property have also $G_{33}$ *and* $G_{34}$ . indeed if<br>$G_{32} < (\widehat{M}_{32})^{(6)}$ it follows $\frac{dG_{33}}{dt} \leq \left((\widehat{M}_{32})^{(6)}\right)_1 - (a_{33}')^{(6)} G_{33}$ and by integrating<br>$G_{33} \leq \left((\widehat{M}_{32})^{(6)}\right)_2 = G_{33}^0 + 2(a_{33})^{(6)}\left((\widehat{M}_{32})^{(6)}\right)_1 / (a_{33}')^{(6)}$<br>In the same way , one can obtain<br>$G_{34} \leq \left((\widehat{M}_{32})^{(6)}\right)_3 = G_{34}^0 + 2(a_{34})^{(6)}\left((\widehat{M}_{32})^{(6)}\right)_2 / (a_{34}')^{(6)}$ | 251 |

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| If $G_{33}$ *or* $G_{34}$ is bounded, the same property follows for $G_{32}$ , $G_{34}$ and $G_{32}$ , $G_{33}$ respectively. | |
| **Remark 29:** If $G_{32}$ *is* bounded, from below, the same property holds for $G_{33}$ *and* $G_{34}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{33}$ is bounded from below. | 252 |
| **Remark 30:** If $T_{32}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(6)}((G_{35})(t),t)) = (b_{33}^{'})^{(6)}$ then $T_{33} \to \infty$.<br><br>**Definition of** $(m)^{(6)}$ and $\varepsilon_6$ :<br><br>Indeed let $t_6$ be so that for $t > t_6$<br><br>$(b_{33})^{(6)} - (b_i^{''})^{(6)}\big((G_{35})(t),t\big) < \varepsilon_6, T_{32}(t) > (m)^{(6)}$ | 253 |
| Then $\frac{dT_{33}}{dt} \geq (a_{33})^{(6)}(m)^{(6)} - \varepsilon_6 T_{33}$ which leads to<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{\varepsilon_6}\right)(1 - e^{-\varepsilon_6 t}) + T_{33}^0 e^{-\varepsilon_6 t}$ If we take t such that $e^{-\varepsilon_6 t} = \frac{1}{2}$ it results<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{2}\right)$, $t = log\frac{2}{\varepsilon_6}$ By taking now $\varepsilon_6$ sufficiently small one sees that $T_{33}$ is unbounded.<br>The same property holds for $T_{34}$ if $\lim_{t\to\infty}(b_{34}^{''})^{(6)}\left((G_{35})(t),t(t),t\right) = (b_{34}^{'})^{(6)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 254 |
| Analogous inequalities hold also for $G_{37}$ , $G_{38}, T_{36}, T_{37}, T_{38}$<br><br>It is now sufficient to take $\frac{(a_i)^{(7)}}{(\hat{M}_{36})^{(7)}}$ , $\frac{(b_i)^{(7)}}{(\hat{M}_{36})^{(7)}} < 1$ and to choose<br>$(\hat{P}_{36})^{(7)}$ and $(\hat{Q}_{36})^{(7)}$ large to have | 255 |
| $\frac{(a_i)^{(7)}}{(\hat{M}_{36})^{(7)}}\left[(\hat{P}_{36})^{(7)} + \left((\hat{P}_{36})^{(7)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{36})^{(7)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{36})^{(7)}$ | 256 |
| $\frac{(b_i)^{(7)}}{(\hat{M}_{36})^{(7)}}\left[\left((\hat{Q}_{36})^{(7)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{36})^{(7)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{36})^{(7)}\right] \leq (\hat{Q}_{36})^{(7)}$ | 257 |
| In order that the operator $\mathcal{A}^{(7)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(7)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{39})^{(1)},(T_{39})^{(1)}),((G_{39})^{(2)},(T_{39})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\hat{M}_{36})^{(7)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\hat{M}_{36})^{(7)}t}\}$<br><br>Indeed if we denote | 258 |

**Definition of** $\widetilde{(G_{39})}, \widetilde{(T_{39})} : \left( \widetilde{(G_{39})}, \widetilde{(T_{39})} \right) = \mathcal{A}^{(7)}((G_{39}), (T_{39}))$
It results

$\left|\tilde{G}_{36}^{(1)} - \tilde{G}_{i}^{(2)}\right| \leq \int_0^t (a_{36})^{(7)} \left|G_{37}^{(1)} - G_{37}^{(2)}\right| e^{-(\widehat{M}_{36})^{(7)} s_{(36)}} e^{(\widehat{M}_{36})^{(7)} s_{(36)}}\, ds_{(36)} +$
$\int_0^t \{(a_{36}^{'})^{(7)} \left|G_{36}^{(1)} - G_{36}^{(2)}\right| e^{-(\widehat{M}_{36})^{(7)} s_{(36)}} e^{-(\widehat{M}_{36})^{(7)} s_{(36)}} +$
$(a_{36}^{''})^{(7)}\left(T_{37}^{(1)}, s_{(36)}\right) \left|G_{36}^{(1)} - G_{36}^{(2)}\right| e^{-(\widehat{M}_{36})^{(7)} s_{(36)}} e^{(\widehat{M}_{36})^{(7)} s_{(36)}} +$
$G_{36}^{(2)} \left|(a_{36}^{''})^{(7)}\left(T_{37}^{(1)}, s_{(36)}\right) - (a_{36}^{''})^{(7)}\left(T_{37}^{(2)}, s_{(36)}\right)\right| \; e^{-(\widehat{M}_{36})^{(7)} s_{(36)}} e^{(\widehat{M}_{36})^{(7)} s_{(36)}} \} ds_{(36)}$

Where $s_{(36)}$ represents integrand that is integrated over the interval $[0, t]$

From the hypotheses on it follows

$\left|(G_{39})^{(1)} - (G_{39})^{(2)}\right| e^{-(\widehat{M}_{36})^{(7)} t} \leq$
$\frac{1}{(\widehat{M}_{36})^{(7)}} \left((a_{36})^{(7)} + (a_{36}^{'})^{(7)} + (\widehat{A}_{36})^{(7)} + (\widehat{P}_{36})^{(7)} (\widehat{k}_{36})^{(7)}\right) d\left(((G_{39})^{(1)}, (T_{39})^{(1)};\ (G_{39})^{(2)}, (T_{39})^{(2)})\right)$

And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows

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**Remark 31:** The fact that we supposed $(a_{36}^{''})^{(7)}$ and $(b_{36}^{''})^{(7)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{36})^{(7)} e^{(\widehat{M}_{36})^{(7)} t}$ *and* $(\widehat{Q}_{36})^{(7)} e^{(\widehat{M}_{36})^{(7)} t}$ respectively of $\mathbb{R}_+$.

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(7)}$ and $(b_i^{''})^{(7)}, i = 36,37,38$ depend only on $T_{37}$ and respectively on $(G_{39})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition.

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**Remark 32:** There does not exist any $t$ where $G_i\,(t) = 0$ *and* $T_i\,(t) = 0$

it results

$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(7)} - (a_i^{''})^{(7)} (T_{37}(s_{(36)}), s_{(36)})\} ds_{(36)}\right]} \geq 0$

$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(7)} t\right)} > 0$ for t > 0

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**Definition of** $\left((\widehat{M}_{36})^{(7)}\right)_1, \left((\widehat{M}_{36})^{(7)}\right)_2$ *and* $\left((\widehat{M}_{36})^{(7)}\right)_3$ :

**Remark 33:** if $G_{36}$ is bounded, the same property have also $G_{37}$ *and* $G_{38}$ . indeed if

$G_{36} < (\widehat{M}_{36})^{(7)}$ it follows $\frac{dG_{37}}{dt} \leq \left((\widehat{M}_{36})^{(7)}\right)_1 - (a_{37}^{'})^{(7)} G_{37}$ and by integrating

$G_{37} \leq \left((\widehat{M}_{36})^{(7)}\right)_2 = G_{37}^0 + 2(a_{37})^{(7)} \left((\widehat{M}_{36})^{(7)}\right)_1 / (a_{37}^{'})^{(7)}$

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| In the same way , one can obtain<br>$G_{38} \leq \left((\widehat{M}_{36})^{(7)}\right)_3 = G_{38}^0 + 2(a_{38})^{(7)}\left((\widehat{M}_{36})^{(7)}\right)_2/(a_{38}')^{(7)}$<br>If $G_{37}$ *or* $G_{38}$ is bounded, the same property follows for $G_{36}$ , $G_{38}$ and $G_{36}$ , $G_{37}$ respectively. | |
| **Remark 34:** If $G_{36}$ *is* bounded, from below, the same property holds for $G_{37}$ *and* $G_{38}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{37}$ is bounded from below. | 263 |
| **Remark 35:** If $T_{36}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(7)}((G_{39})(t),t)) = (b_{37}')^{(7)}$ then $T_{37} \to \infty$.<br>**Definition of** $(m)^{(7)}$ and $\varepsilon_7$ :<br>Indeed let $t_7$ be so that for $t > t_7$<br>$(b_{37})^{(7)} - (b_i'')^{(7)}((G_{39})(t),t) < \varepsilon_7, T_{36}(t) > (m)^{(7)}$ | 264 |
| Then $\frac{dT_{37}}{dt} \geq (a_{37})^{(7)}(m)^{(7)} - \varepsilon_7 T_{37}$ which leads to<br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{\varepsilon_7}\right)(1 - e^{-\varepsilon_7 t}) + T_{37}^0 e^{-\varepsilon_7 t}$ If we take t such that $e^{-\varepsilon_7 t} = \frac{1}{2}$ it results<br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{2}\right)$, $t = log\frac{2}{\varepsilon_7}$ By taking now $\varepsilon_7$ sufficiently small one sees that $T_{37}$ is unbounded.<br>The same property holds for $T_{38}$ if $\lim_{t\to\infty}(b_{38}'')^{(7)}\left((G_{39})(t),t\right) = (b_{38}')^{(7)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 265 |
| It is now sufficient to take $\frac{(a_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}$ , $\frac{(b_i)^{(8)}}{(\widehat{M}_{40})^{(8)}} < 1$ and to choose<br>$(\widehat{P}_{40})^{(8)}$ and $(\widehat{Q}_{40})^{(8)}$ large to have | 266 |
| $\frac{(a_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}\left[(\widehat{P}_{40})^{(8)} + \left((\widehat{P}_{40})^{(8)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{40})^{(8)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{40})^{(8)}$ | 267 |
| $\frac{(b_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}\left[\left((\widehat{Q}_{40})^{(8)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{40})^{(8)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{40})^{(8)}\right] \leq (\widehat{Q}_{40})^{(8)}$ | 268 |
| In order that the operator $\mathcal{A}^{(8)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(8)}$ is a contraction with respect to the metric | |

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| $d\left(((G_{43})^{(1)},(T_{43})^{(1)}),((G_{43})^{(2)},(T_{43})^{(2)})\right) =$ $\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{40})^{(8)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{40})^{(8)}t}\}$ | 269 |
| Indeed if we denote<br><br>**Definition of** $\widetilde{(G_{43})},\widetilde{(T_{43})}$ : $\left(\widetilde{(G_{43})},\widetilde{(T_{43})}\right)=\mathcal{A}^{(8)}((G_{43}),(T_{43}))$ | 270 |
| It results<br><br>$\left\|\tilde{G}_{40}^{(1)}-\tilde{G}_i^{(2)}\right\|\le\int_0^t(a_{40})^{(8)}\left\|G_{41}^{(1)}-G_{41}^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{(\widehat{M}_{40})^{(8)}s_{(40)}}\,ds_{(40)}+$<br>$\int_0^t\{(a_{40}^{'})^{(8)}\left\|G_{40}^{(1)}-G_{40}^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}+$<br>$(a_{40}^{''})^{(8)}\left(T_{41}^{(1)},s_{(40)}\right)\left\|G_{40}^{(1)}-G_{40}^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{(\widehat{M}_{40})^{(8)}s_{(40)}}+$<br>$G_{40}^{(2)}\left\\|(a_{40}^{''})^{(8)}\left(T_{41}^{(1)},s_{(40)}\right)-(a_{40}^{''})^{(8)}\left(T_{41}^{(2)},s_{(40)}\right)\right\\|\ e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{(\widehat{M}_{40})^{(8)}s_{(40)}}\}ds_{(40)}$ | 271 |
| Where $s_{(40)}$ represents integrand that is integrated over the interval $[0,\mathrm{t}]$<br><br>From the hypotheses it follows | 272 |
| $\left\|(G_{43})^{(1)}-(G_{43})^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}t}\le$<br>$\frac{1}{(\widehat{M}_{40})^{(8)}}\left((a_{40})^{(8)}+(a_{40}^{'})^{(8)}+(\widehat{A}_{40})^{(8)}+(\widehat{P}_{40})^{(8)}(\widehat{k}_{40})^{(8)}\right)d\left(((G_{43})^{(1)},(T_{43})^{(1)};\ (G_{43})^{(2)},(T_{43})^{(2)})\right)$<br><br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 273 |
| **Remark 36:** The fact that we supposed $(a_{40}^{''})^{(8)}$ and $(b_{40}^{''})^{(8)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{40})^{(8)}e^{(\widehat{M}_{40})^{(8)}t}$ *and* $(\widehat{Q}_{40})^{(8)}e^{(\widehat{M}_{40})^{(8)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(8)}$ and $(b_i^{''})^{(8)}, i=40,41,42$ depend only on $\mathrm{T}_{41}$ and respectively on $(G_{43})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 274 |
| **Remark 37** There does not exist any $t$ where $G_i\,(t)=0$ *and* $T_i\,(t)=0$<br><br>it results<br><br>$G_i\,(t)\ge G_i^0e^{\left[-\int_0^t\{(a_i^{'})^{(8)}-(a_i^{''})^{(8)}(T_{41}(s_{(40)}),s_{(40)})\}ds_{(40)}\right]}\ge 0$<br><br>$T_i\,(t)\ge T_i^0e^{\left(-(b_i^{'})^{(8)}t\right)}>0$ for t > 0 | 275 |
| **Definition of** $\left((\widehat{M}_{40})^{(8)}\right)_1,\left((\widehat{M}_{40})^{(8)}\right)_2$ *and* $\left((\widehat{M}_{40})^{(8)}\right)_3$ : | 276 |

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| **Remark 38:** if $G_{40}$ is bounded, the same property have also $G_{41}\ and\ G_{42}$ . indeed if<br>$G_{40} < (\widehat{M}_{40})^{(8)}$ it follows $\frac{dG_{41}}{dt} \leq \left((\widehat{M}_{40})^{(8)}\right)_1 - (a'_{41})^{(8)}G_{41}$ and by integrating<br>$G_{41} \leq \left((\widehat{M}_{40})^{(8)}\right)_2 = G_{41}^0 + 2(a_{41})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_1/(a'_{41})^{(8)}$<br>In the same way , one can obtain<br>$G_{42} \leq \left((\widehat{M}_{40})^{(8)}\right)_3 = G_{42}^0 + 2(a_{42})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_2/(a'_{42})^{(8)}$<br>If $G_{41}\ or\ G_{42}$ is bounded, the same property follows for $G_{40}$ , $G_{42}$ and $G_{40}$ , $G_{41}$ respectively. | |
| **Remark 39:** If $G_{40}$ *is* bounded, from below, the same property holds for $G_{41}\ and\ G_{42}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{41}$ is bounded from below. | 277 |
| **Remark 40:** If $T_{40}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(8)}((G_{43})(t),t)) = (b'_{41})^{(8)}$ then $T_{41} \to \infty$.<br>**Definition of** $(m)^{(8)}$ and $\varepsilon_8$ :<br>Indeed let $t_8$ be so that for $t > t_8$<br>$(b_{41})^{(8)} - (b_i^{''})^{(8)}\left((G_{43})(t),t\right) < \varepsilon_8, T_{40}(t) > (m)^{(8)}$ | 278 |
| Then $\frac{dT_{41}}{dt} \geq (a_{41})^{(8)}(m)^{(8)} - \varepsilon_8 T_{41}$ which leads to<br>$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{\varepsilon_8}\right)(1 - e^{-\varepsilon_8 t}) + T_{41}^0 e^{-\varepsilon_8 t}$ If we take t such that $e^{-\varepsilon_8 t} = \frac{1}{2}$ it results<br>$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{2}\right),\ t = log\frac{2}{\varepsilon_8}$ By taking now $\varepsilon_8$ sufficiently small one sees that $T_{41}$ is unbounded.<br>The same property holds for $T_{42}$ if $\lim_{t\to\infty}(b_{42}^{''})^{(8)}\left((G_{43})(t),t(t),t\right) = (b'_{42})^{(8)}$ | 279 |
| It is now sufficient to take $\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}$ , $\frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}} < 1$ and to choose $(\widehat{P}_{44})^{(9)}$ and $(\widehat{Q}_{44})^{(9)}$ large to have | 279 A |
| $\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[(\widehat{P}_{44})^{(9)} + \left((\widehat{P}_{44})^{(9)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{44})^{(9)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{44})^{(9)}$ | |
| $\frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[\left((\widehat{Q}_{44})^{(9)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{44})^{(9)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{44})^{(9)}\right] \leq (\widehat{Q}_{44})^{(9)}$ | |

In order that the operator $\mathcal{A}^{(9)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying 39,35,36 into itself

The operator $\mathcal{A}^{(9)}$ is a contraction with respect to the metric

$$d\left(((G_{47})^{(1)},(T_{47})^{(1)}),((G_{47})^{(2)},(T_{47})^{(2)})\right)=$$
$$sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t)-G_i^{(2)}(t)\right|e^{-(\widehat{M}_{44})^{(9)}t},\max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t)-T_i^{(2)}(t)\right|e^{-(\widehat{M}_{44})^{(9)}t}\}$$

Indeed if we denote

**Definition of** $\widetilde{(G_{47})},\widetilde{(T_{47})}:\left(\widetilde{(G_{47})},\widetilde{(T_{47})}\right)=\mathcal{A}^{(9)}\left((G_{47}),(T_{47})\right)$
It results

$$\left|\tilde{G}_{44}^{(1)}-\tilde{G}_i^{(2)}\right|\le\int_0^t(a_{44})^{(9)}\left|G_{45}^{(1)}-G_{45}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\,ds_{(44)}+$$
$$\int_0^t\{(a_{44}^{'})^{(9)}\left|G_{44}^{(1)}-G_{44}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}+$$
$$(a_{44}^{''})^{(9)}\left(T_{45}^{(1)},s_{(44)}\right)\left|G_{44}^{(1)}-G_{44}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}+$$
$$G_{44}^{(2)}\left|(a_{44}^{''})^{(9)}\left(T_{45}^{(1)},s_{(44)}\right)-(a_{44}^{''})^{(9)}\left(T_{45}^{(2)},s_{(44)}\right)\right|\ e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\}ds_{(44)}$$

Where $s_{(44)}$ represents integrand that is integrated over the interval $[0,\mathrm{t}]$

From the hypotheses on 45,46,47,28 and 29 it follows

$$\left|(G_{47})^{(1)}-G^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}t}\le$$
$$\frac{1}{(\widehat{M}_{44})^{(9)}}\left((a_{44})^{(9)}+(a_{44}^{'})^{(9)}+(\widehat{A}_{44})^{(9)}+(\widehat{P}_{44})^{(9)}(\widehat{k}_{44})^{(9)}\right)d\left(((G_{47})^{(1)},(T_{47})^{(1)};\ (G_{47})^{(2)},(T_{47})^{(2)})\right)$$

And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis (39,35,36) the result follows

**Remark 41:** The fact that we supposed $(a_{44}^{''})^{(9)}$ and $(b_{44}^{''})^{(9)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{44})^{(9)}e^{(\widehat{M}_{44})^{(9)}t}$ *and* $(\widehat{Q}_{44})^{(9)}e^{(\widehat{M}_{44})^{(9)}t}$ respectively of $\mathbb{R}_+$.

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(9)}$ and $(b_i^{''})^{(9)}, i=44{,}45{,}46$ depend only on $\mathrm{T}_{45}$ and respectively on $(G_{47})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition.

**Remark 42:** There does not exist any $t$ where $G_i\,(t)=0$ *and* $T_i\,(t)=0$

From 99 to 44 it results

$$G_i\,(t)\ge G_i^0e^{\left[-\int_0^t\{(a_i^{'})^{(9)}-(a_i^{''})^{(9)}(T_{45}(s_{(44)}),s_{(44)})\}ds_{(44)}\right]}\ge 0$$

$$T_i\,(t)\ge T_i^0e^{(-(b_i^{'})^{(9)}t)}>0\quad\text{for t}>0$$

**Definition of** $\left((\widehat{M}_{44})^{(9)}\right)_1,\left((\widehat{M}_{44})^{(9)}\right)_2$ *and* $\left((\widehat{M}_{44})^{(9)}\right)_3$ :

**Remark 43:** if $G_{44}$ is bounded, the same property have also $G_{45}$ *and* $G_{46}$ . indeed if

$G_{44} < (\widehat{M}_{44})^{(9)}$ it follows $\frac{dG_{45}}{dt} \leq \left((\widehat{M}_{44})^{(9)}\right)_1 - (a_{45}^{'})^{(9)} G_{45}$ and by integrating

$G_{45} \leq \left((\widehat{M}_{44})^{(9)}\right)_2 = G_{45}^0 + 2(a_{45})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_1/(a_{45}^{'})^{(9)}$

In the same way , one can obtain

$G_{46} \leq \left((\widehat{M}_{44})^{(9)}\right)_3 = G_{46}^0 + 2(a_{46})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_2/(a_{46}^{'})^{(9)}$

If $G_{45}$ *or* $G_{46}$ is bounded, the same property follows for $G_{44}$ , $G_{46}$ and $G_{44}$ , $G_{45}$ respectively.

**Remark 44:** If $G_{44}$ *is* bounded, from below, the same property holds for $G_{45}$ *and* $G_{46}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{45}$ is bounded from below.

**Remark 45:** If $T_{44}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(9)}\left((G_{47})(t),t\right) = (b_{45}^{'})^{(9)}$ then $T_{45} \to \infty$.

**Definition of** $(m)^{(9)}$ and $\varepsilon_9$ :

Indeed let $t_9$ be so that for $t > t_9$

$$(b_{45})^{(9)} - (b_i^{''})^{(9)}\left((G_{47})(t),t\right) < \varepsilon_9, T_{44}(t) > (m)^{(9)}$$

Then $\frac{dT_{45}}{dt} \geq (a_{45})^{(9)}(m)^{(9)} - \varepsilon_9 T_{45}$ which leads to

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{\varepsilon_9}\right)(1 - e^{-\varepsilon_9 t}) + T_{45}^0 e^{-\varepsilon_9 t}$ If we take t such that $e^{-\varepsilon_9 t} = \frac{1}{2}$ it results

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{2}\right)$, $t = log\frac{2}{\varepsilon_9}$ By taking now $\varepsilon_9$ sufficiently small one sees that $T_{45}$ is unbounded.
The same property holds for $T_{46}$ if $\lim_{t\to\infty}(b_{46}^{''})^{(9)}\left((G_{47})(t),t\right) = (b_{46}^{'})^{(9)}$
We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 92

**Behavior of the solutions of equation** 280

**Theorem** If we denote and define

**Definition of** $(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ :

$(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ four constants satisfying
$-(\sigma_2)^{(1)} \leq -(a_{13}^{'})^{(1)} + (a_{14}^{'})^{(1)} - (a_{13}^{''})^{(1)}(T_{14},t) + (a_{14}^{''})^{(1)}(T_{14},t) \leq -(\sigma_1)^{(1)}$

$-(\tau_2)^{(1)} \leq -(b_{13}^{'})^{(1)} + (b_{14}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G,t) - (b_{14}^{''})^{(1)}(G,t) \leq -(\tau_1)^{(1)}$

**Definition of** $(v_1)^{(1)}, (v_2)^{(1)}, (u_1)^{(1)}, (u_2)^{(1)}, v^{(1)}, u^{(1)}$ : 281

By $(v_1)^{(1)} > 0$ , $(v_2)^{(1)} < 0$ and respectively $(u_1)^{(1)} > 0$ , $(u_2)^{(1)} < 0$ the roots of the equations
$(a_{14})^{(1)}\left(v^{(1)}\right)^2 + (\sigma_1)^{(1)}v^{(1)} - (a_{13})^{(1)} = 0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2 + (\tau_1)^{(1)}u^{(1)} - (b_{13})^{(1)} = 0$

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| **Definition of** $(\bar{v}_1)^{(1)},,(\bar{v}_2)^{(1)},(\bar{u}_1)^{(1)},(\bar{u}_2)^{(1)}$ :<br><br>By $(\bar{v}_1)^{(1)}>0\,,(\bar{v}_2)^{(1)}<0$ and respectively $(\bar{u}_1)^{(1)}>0\,,(\bar{u}_2)^{(1)}<0$ the roots of the equations $(a_{14})^{(1)}\left(v^{(1)}\right)^2+(\sigma_2)^{(1)}v^{(1)}-(a_{13})^{(1)}=0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2+(\tau_2)^{(1)}u^{(1)}-(b_{13})^{(1)}=0$ | 282 |
| **Definition of** $(m_1)^{(1)}\,,(m_2)^{(1)}\,,(\mu_1)^{(1)},(\mu_2)^{(1)},(\nu_0)^{(1)}$ :-<br><br>If we define $(m_1)^{(1)}\,,(m_2)^{(1)}\,,(\mu_1)^{(1)},(\mu_2)^{(1)}$ by<br>$(m_2)^{(1)}=(\nu_0)^{(1)},(m_1)^{(1)}=(\nu_1)^{(1)},\ if\ (\nu_0)^{(1)}<(\nu_1)^{(1)}$<br><br>$(m_2)^{(1)}=(\nu_1)^{(1)},(m_1)^{(1)}=(\bar{\nu}_1)^{(1)}\,,if\ (\nu_1)^{(1)}<(\nu_0)^{(1)}<(\bar{\nu}_1)^{(1)},$<br><br>and $\boxed{(\nu_0)^{(1)}=\frac{G_{13}^0}{G_{14}^0}}$<br><br>$(\,m_2)^{(1)}=(\nu_1)^{(1)},(m_1)^{(1)}=(\nu_0)^{(1)},\ if\ (\bar{\nu}_1)^{(1)}<(\nu_0)^{(1)}$ | 283 |
| and analogously<br><br>$(\mu_2)^{(1)}=(u_0)^{(1)},(\mu_1)^{(1)}=(u_1)^{(1)},\ if\ (u_0)^{(1)}<(u_1)^{(1)}$<br><br>$(\mu_2)^{(1)}=(u_1)^{(1)},(\mu_1)^{(1)}=(\bar{u}_1)^{(1)}\,,if\ (u_1)^{(1)}<(u_0)^{(1)}<(\bar{u}_1)^{(1)},$<br><br>and $\boxed{(u_0)^{(1)}=\frac{T_{13}^0}{T_{14}^0}}$<br><br>$(\,\mu_2)^{(1)}=(u_1)^{(1)},(\mu_1)^{(1)}=(u_0)^{(1)},if\ (\bar{u}_1)^{(1)}<(u_0)^{(1)}$ where $(u_1)^{(1)},(\bar{u}_1)^{(1)}$<br><br>are defined | 284 |
| Then the solution of global equations satisfies the inequalities<br><br>$G_{13}^0e^{((S_1)^{(1)}-(p_{13})^{(1)})t}\le G_{13}(t)\le G_{13}^0e^{(S_1)^{(1)}t}$<br><br>where $(p_i)^{(1)}$ is defined by equation<br><br>$\frac{1}{(m_1)^{(1)}}G_{13}^0e^{((S_1)^{(1)}-(p_{13})^{(1)})t}\le G_{14}(t)\le\frac{1}{(m_2)^{(1)}}G_{13}^0e^{(S_1)^{(1)}t}$ | 285 |
| $(\frac{(a_{15})^{(1)}G_{13}^0}{(m_1)^{(1)}((S_1)^{(1)}-(p_{13})^{(1)}-(S_2)^{(1)})}\left[e^{((S_1)^{(1)}-(p_{13})^{(1)})t}-e^{-(S_2)^{(1)}t}\right]+G_{15}^0e^{-(S_2)^{(1)}t}\le G_{15}(t)\le$<br><br>$\frac{(a_{15})^{(1)}G_{13}^0}{(m_2)^{(1)}((S_1)^{(1)}-(a_{15}^{'})^{(1)})}[e^{(S_1)^{(1)}t}-e^{-(a_{15}^{'})^{(1)}t}]+\ G_{15}^0e^{-(a_{15}^{'})^{(1)}t})$ | 286 |
| $\boxed{T_{13}^0e^{(R_1)^{(1)}t}\le T_{13}(t)\le T_{13}^0e^{((R_1)^{(1)}+(r_{13})^{(1)})t}}$ | 287 |
| $\frac{1}{(\mu_1)^{(1)}}T_{13}^0e^{(R_1)^{(1)}t}\le T_{13}(t)\le\frac{1}{(\mu_2)^{(1)}}T_{13}^0e^{((R_1)^{(1)}+(r_{13})^{(1)})t}$ | 288 |

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| $\frac{(b_{15})^{(1)}T_{13}^{0}}{(\mu_1)^{(1)}((R_1)^{(1)}-(b_{15}^{'})^{(1)})}\left[e^{(R_1)^{(1)}t}-e^{-(b_{15}^{'})^{(1)}t}\right]+T_{15}^{0}e^{-(b_{15}^{'})^{(1)}t}\leq T_{15}(t)\leq$ $\frac{(a_{15})^{(1)}T_{13}^{0}}{(\mu_2)^{(1)}((R_1)^{(1)}+(r_{13})^{(1)}+(R_2)^{(1)})}\left[e^{((R_1)^{(1)}+(r_{13})^{(1)})t}-e^{-(R_2)^{(1)}t}\right]+T_{15}^{0}e^{-(R_2)^{(1)}t}$ | 289 |
| **Definition of** $(S_1)^{(1)},(S_2)^{(1)},(R_1)^{(1)},(R_2)^{(1)}$:- Where $(S_1)^{(1)}=(a_{13})^{(1)}(m_2)^{(1)}-(a_{13}^{'})^{(1)}$ $(S_2)^{(1)}=(a_{15})^{(1)}-(p_{15})^{(1)}$ $(R_1)^{(1)}=(b_{13})^{(1)}(\mu_2)^{(1)}-(b_{13}^{'})^{(1)}$ $(R_2)^{(1)}=(b_{15}^{'})^{(1)}-(r_{15})^{(1)}$ | 290 |
| **Behavior of the solutions of equation** **Theorem 2:** If we denote and define | 291 |
| **Definition of** $(\sigma_1)^{(2)},(\sigma_2)^{(2)},(\tau_1)^{(2)},(\tau_2)^{(2)}$ : $(\sigma_1)^{(2)},(\sigma_2)^{(2)},(\tau_1)^{(2)},(\tau_2)^{(2)}$ four constants satisfying | 292 |
| $-(\sigma_2)^{(2)}\leq-(a_{16}^{'})^{(2)}+(a_{17}^{'})^{(2)}-(a_{16}^{''})^{(2)}(\mathrm{T}_{17},t)+(a_{17}^{''})^{(2)}(\mathrm{T}_{17},t)\leq-(\sigma_1)^{(2)}$ | 293 |
| $-(\tau_2)^{(2)}\leq-(b_{16}^{'})^{(2)}+(b_{17}^{'})^{(2)}-(b_{16}^{''})^{(2)}\big((G_{19}),t\big)-(b_{17}^{''})^{(2)}\big((G_{19}),t\big)\leq-(\tau_1)^{(2)}$ | 294 |
| **Definition of** $(v_1)^{(2)},(v_2)^{(2)},(u_1)^{(2)},(u_2)^{(2)}$ : | 295 |
| By $(v_1)^{(2)}>0$ , $(v_2)^{(2)}<0$ and respectively $(u_1)^{(2)}>0$ , $(u_2)^{(2)}<0$ the roots | 296 |
| of the equations $(a_{17})^{(2)}\big(v^{(2)}\big)^2+(\sigma_1)^{(2)}v^{(2)}-(a_{16})^{(2)}=0$ | 297 |
| and $(b_{14})^{(2)}\big(u^{(2)}\big)^2+(\tau_1)^{(2)}u^{(2)}-(b_{16})^{(2)}=0$ and | 298 |
| **Definition of** $(\bar{v}_1)^{(2)},,(\bar{v}_2)^{(2)},(\bar{u}_1)^{(2)},(\bar{u}_2)^{(2)}$ : | 299 |
| By $(\bar{v}_1)^{(2)}>0$ , $(\bar{v}_2)^{(2)}<0$ and respectively $(\bar{u}_1)^{(2)}>0$ , $(\bar{u}_2)^{(2)}<0$ the | 300 |
| roots of the equations $(a_{17})^{(2)}\big(v^{(2)}\big)^2+(\sigma_2)^{(2)}v^{(2)}-(a_{16})^{(2)}=0$ | 301 |
| and $(b_{17})^{(2)}\big(u^{(2)}\big)^2+(\tau_2)^{(2)}u^{(2)}-(b_{16})^{(2)}=0$ | 302 |
| **Definition of** $(m_1)^{(2)},(m_2)^{(2)},(\mu_1)^{(2)},(\mu_2)^{(2)}$ :- | 303 |
| If we define $(m_1)^{(2)},(m_2)^{(2)},(\mu_1)^{(2)},(\mu_2)^{(2)}$ by | 304 |
| $(m_2)^{(2)}=(v_0)^{(2)},(m_1)^{(2)}=(v_1)^{(2)}$, ***if*** $(v_0)^{(2)}<(v_1)^{(2)}$ | 305 |

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| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ , $\boldsymbol{if}$ $(\nu_1)^{(2)} < (\nu_0)^{(2)} < (\bar{\nu}_1)^{(2)}$, and $\boxed{(\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0}}$ | 306 |
| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\nu_0)^{(2)}$, $\boldsymbol{if}$ $(\bar{\nu}_1)^{(2)} < (\nu_0)^{(2)}$ | 307 |
| and analogously<br>$(\mu_2)^{(2)} = (u_0)^{(2)}, (\mu_1)^{(2)} = (u_1)^{(2)}$, $\boldsymbol{if}$ $(u_0)^{(2)} < (u_1)^{(2)}$<br>$(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (\bar{u}_1)^{(2)}$ , $\boldsymbol{if}$ $(u_1)^{(2)} < (u_0)^{(2)} < (\bar{u}_1)^{(2)}$,<br>and $\boxed{(u_0)^{(2)} = \frac{T_{16}^0}{T_{17}^0}}$ | 308 |
| $(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (u_0)^{(2)}$, $\boldsymbol{if}$ $(\bar{u}_1)^{(2)} < (u_0)^{(2)}$ | 309 |
| Then the solution of global equations satisfies the inequalities<br>$G_{16}^0 e^{((S_1)^{(2)} - (p_{16})^{(2)})t} \le G_{16}(t) \le G_{16}^0 e^{(S_1)^{(2)} t}$ | 310 |
| $(p_i)^{(2)}$ is defined by equation | |
| $\frac{1}{(m_1)^{(2)}} G_{16}^0 e^{((S_1)^{(2)} - (p_{16})^{(2)})t} \le G_{17}(t) \le \frac{1}{(m_2)^{(2)}} G_{16}^0 e^{(S_1)^{(2)} t}$ | 311 |
| $\left(\frac{(a_{18})^{(2)} G_{16}^0}{(m_1)^{(2)}((S_1)^{(2)} - (p_{16})^{(2)} - (S_2)^{(2)})}\left[e^{((S_1)^{(2)} - (p_{16})^{(2)})t} - e^{-(S_2)^{(2)} t}\right] + G_{18}^0 e^{-(S_2)^{(2)} t} \le G_{18}(t) \le \right.$<br>$\left.\frac{(a_{18})^{(2)} G_{16}^0}{(m_2)^{(2)}((S_1)^{(2)} - (a'_{18})^{(2)})}\left[e^{(S_1)^{(2)} t} - e^{-(a'_{18})^{(2)} t}\right] + G_{18}^0 e^{-(a'_{18})^{(2)} t}\right)$ | 312 |
| $\boxed{T_{16}^0 e^{(R_1)^{(2)} t} \le T_{16}(t) \le T_{16}^0 e^{((R_1)^{(2)} + (r_{16})^{(2)})t}}$ | 313 |
| $\frac{1}{(\mu_1)^{(2)}} T_{16}^0 e^{(R_1)^{(2)} t} \le T_{16}(t) \le \frac{1}{(\mu_2)^{(2)}} T_{16}^0 e^{((R_1)^{(2)} + (r_{16})^{(2)})t}$ | 314 |
| $\frac{(b_{18})^{(2)} T_{16}^0}{(\mu_1)^{(2)}((R_1)^{(2)} - (b'_{18})^{(2)})}\left[e^{(R_1)^{(2)} t} - e^{-(b'_{18})^{(2)} t}\right] + T_{18}^0 e^{-(b'_{18})^{(2)} t} \le T_{18}(t) \le$<br>$\frac{(a_{18})^{(2)} T_{16}^0}{(\mu_2)^{(2)}((R_1)^{(2)} + (r_{16})^{(2)} + (R_2)^{(2)})}\left[e^{((R_1)^{(2)} + (r_{16})^{(2)})t} - e^{-(R_2)^{(2)} t}\right] + T_{18}^0 e^{-(R_2)^{(2)} t}$ | 315 |
| **Definition of** $(S_1)^{(2)}, (S_2)^{(2)}, (R_1)^{(2)}, (R_2)^{(2)}$:- | 316 |
| Where $(S_1)^{(2)} = (a_{16})^{(2)}(m_2)^{(2)} - (a'_{16})^{(2)}$<br>$(S_2)^{(2)} = (a_{18})^{(2)} - (p_{18})^{(2)}$ | 317 |
| $(R_1)^{(2)} = (b_{16})^{(2)}(\mu_2)^{(1)} - (b'_{16})^{(2)}$ | 318 |

$(R_2)^{(2)} = (b_{18}')^{(2)} - (r_{18})^{(2)}$

**Behavior of the solutions** 319

**Theorem 3:** If we denote and define

**Definition of** $(\sigma_1)^{(3)}, (\sigma_2)^{(3)}, (\tau_1)^{(3)}, (\tau_2)^{(3)}$ :

$(\sigma_1)^{(3)}, (\sigma_2)^{(3)}, (\tau_1)^{(3)}, (\tau_2)^{(3)}$ four constants satisfying

$$-(\sigma_2)^{(3)} \le -(a_{20}')^{(3)} + (a_{21}')^{(3)} - (a_{20}'')^{(3)}(T_{21}, t) + (a_{21}'')^{(3)}(T_{21}, t) \le -(\sigma_1)^{(3)}$$

$$-(\tau_2)^{(3)} \le -(b_{20}')^{(3)} + (b_{21}')^{(3)} - (b_{20}'')^{(3)}(G_{23}, t) - (b_{21}'')^{(3)}((G_{23}), t) \le -(\tau_1)^{(3)}$$

**Definition of** $(v_1)^{(3)}, (v_2)^{(3)}, (u_1)^{(3)}, (u_2)^{(3)}$ : 320

By $(v_1)^{(3)} > 0$ , $(v_2)^{(3)} < 0$ and respectively $(u_1)^{(3)} > 0$ , $(u_2)^{(3)} < 0$ the roots of the equations $(a_{21})^{(3)}(v^{(3)})^2 + (\sigma_1)^{(3)}v^{(3)} - (a_{20})^{(3)} = 0$

and $(b_{21})^{(3)}(u^{(3)})^2 + (\tau_1)^{(3)}u^{(3)} - (b_{20})^{(3)} = 0$ and

By $(\bar{v}_1)^{(3)} > 0$ , $(\bar{v}_2)^{(3)} < 0$ and respectively $(\bar{u}_1)^{(3)} > 0$ , $(\bar{u}_2)^{(3)} < 0$ the

roots of the equations $(a_{21})^{(3)}(v^{(3)})^2 + (\sigma_2)^{(3)}v^{(3)} - (a_{20})^{(3)} = 0$

and $(b_{21})^{(3)}(u^{(3)})^2 + (\tau_2)^{(3)}u^{(3)} - (b_{20})^{(3)} = 0$

**Definition of** $(m_1)^{(3)}, (m_2)^{(3)}, (\mu_1)^{(3)}, (\mu_2)^{(3)}$ :- 321

If we define $(m_1)^{(3)}, (m_2)^{(3)}, (\mu_1)^{(3)}, (\mu_2)^{(3)}$ by

$(m_2)^{(3)} = (v_0)^{(3)}, (m_1)^{(3)} = (v_1)^{(3)}$, ***if*** $(v_0)^{(3)} < (v_1)^{(3)}$

$(m_2)^{(3)} = (v_1)^{(3)}, (m_1)^{(3)} = (\bar{v}_1)^{(3)}$ , ***if*** $(v_1)^{(3)} < (v_0)^{(3)} < (\bar{v}_1)^{(3)}$,

and $\boxed{(v_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0}}$

$(m_2)^{(3)} = (v_1)^{(3)}, (m_1)^{(3)} = (v_0)^{(3)}$, ***if*** $(\bar{v}_1)^{(3)} < (v_0)^{(3)}$

and analogously 322

$(\mu_2)^{(3)} = (u_0)^{(3)}, (\mu_1)^{(3)} = (u_1)^{(3)}$, ***if*** $(u_0)^{(3)} < (u_1)^{(3)}$

$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (\bar{u}_1)^{(3)}$ , ***if*** $(u_1)^{(3)} < (u_0)^{(3)} < (\bar{u}_1)^{(3)}$, and $\boxed{(u_0)^{(3)} = \frac{T_{20}^0}{T_{21}^0}}$

$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (u_0)^{(3)}$, ***if*** $(\bar{u}_1)^{(3)} < (u_0)^{(3)}$

Then the solution of global equations satisfies the inequalities

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| $G_{20}^0 e^{((S_1)^{(3)}-(p_{20})^{(3)})t} \le G_{20}(t) \le G_{20}^0 e^{(S_1)^{(3)}t}$<br>$(p_i)^{(3)}$ is defined by equation | |
| $\frac{1}{(m_1)^{(3)}} G_{20}^0 e^{((S_1)^{(3)}-(p_{20})^{(3)})t} \le G_{21}(t) \le \frac{1}{(m_2)^{(3)}} G_{20}^0 e^{(S_1)^{(3)}t}$ | 323 |
| $\left(\frac{(a_{22})^{(3)}G_{20}^0}{(m_1)^{(3)}((S_1)^{(3)}-(p_{20})^{(3)}-(S_2)^{(3)})}\left[e^{((S_1)^{(3)}-(p_{20})^{(3)})t} - e^{-(S_2)^{(3)}t}\right] + G_{22}^0 e^{-(S_2)^{(3)}t} \le G_{22}(t) \le \right.$<br>$\left.\frac{(a_{22})^{(3)}G_{20}^0}{(m_2)^{(3)}((S_1)^{(3)}-(a'_{22})^{(3)})}[e^{(S_1)^{(3)}t} - e^{-(a'_{22})^{(3)}t}] + G_{22}^0 e^{-(a'_{22})^{(3)}t}\right)$ | 324 |
| $\boxed{T_{20}^0 e^{(R_1)^{(3)}t} \le T_{20}(t) \le T_{20}^0 e^{((R_1)^{(3)}+(r_{20})^{(3)})t}}$ | 325 |
| $\frac{1}{(\mu_1)^{(3)}} T_{20}^0 e^{(R_1)^{(3)}t} \le T_{20}(t) \le \frac{1}{(\mu_2)^{(3)}} T_{20}^0 e^{((R_1)^{(3)}+(r_{20})^{(3)})t}$ | 326 |
| $\frac{(b_{22})^{(3)}T_{20}^0}{(\mu_1)^{(3)}((R_1)^{(3)}-(b'_{22})^{(3)})}\left[e^{(R_1)^{(3)}t} - e^{-(b'_{22})^{(3)}t}\right] + T_{22}^0 e^{-(b'_{22})^{(3)}t} \le T_{22}(t) \le$<br>$\frac{(a_{22})^{(3)}T_{20}^0}{(\mu_2)^{(3)}((R_1)^{(3)}+(r_{20})^{(3)}+(R_2)^{(3)})}\left[e^{((R_1)^{(3)}+(r_{20})^{(3)})t} - e^{-(R_2)^{(3)}t}\right] + T_{22}^0 e^{-(R_2)^{(3)}t}$ | 327 |
| **Definition of** $(S_1)^{(3)}, (S_2)^{(3)}, (R_1)^{(3)}, (R_2)^{(3)}$:-<br>Where $(S_1)^{(3)} = (a_{20})^{(3)}(m_2)^{(3)} - (a'_{20})^{(3)}$<br>$(S_2)^{(3)} = (a_{22})^{(3)} - (p_{22})^{(3)}$<br>$(R_1)^{(3)} = (b_{20})^{(3)}(\mu_2)^{(3)} - (b'_{20})^{(3)}$<br>$(R_2)^{(3)} = (b'_{22})^{(3)} - (r_{22})^{(3)}$ | 328 |
| **Behavior of the solutions of equation**<br>**Theorem:** If we denote and define<br>**Definition of** $(\sigma_1)^{(4)}, (\sigma_2)^{(4)}, (\tau_1)^{(4)}, (\tau_2)^{(4)}$ :<br>$(\sigma_1)^{(4)}, (\sigma_2)^{(4)}, (\tau_1)^{(4)}, (\tau_2)^{(4)}$ four constants satisfying<br>$-(\sigma_2)^{(4)} \le -(a'_{24})^{(4)} + (a'_{25})^{(4)} - (a''_{24})^{(4)}(T_{25}, t) + (a''_{25})^{(4)}(T_{25}, t) \le -(\sigma_1)^{(4)}$<br>$-(\tau_2)^{(4)} \le -(b'_{24})^{(4)} + (b'_{25})^{(4)} - (b''_{24})^{(4)}((G_{27}), t) - (b''_{25})^{(4)}((G_{27}), t) \le -(\tau_1)^{(4)}$ | |
| **Definition of** $(v_1)^{(4)}, (v_2)^{(4)}, (u_1)^{(4)}, (u_2)^{(4)}, v^{(4)}, u^{(4)}$ :<br>By $(v_1)^{(4)} > 0$, $(v_2)^{(4)} < 0$ and respectively $(u_1)^{(4)} > 0$, $(u_2)^{(4)} < 0$ the roots of the equations $(a_{25})^{(4)}(v^{(4)})^2 + (\sigma_1)^{(4)}v^{(4)} - (a_{24})^{(4)} = 0$<br>and $(b_{25})^{(4)}(u^{(4)})^2 + (\tau_1)^{(4)}u^{(4)} - (b_{24})^{(4)} = 0$ and | 329 |

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| **Definition of** $(\bar{v}_1)^{(4)},,(\bar{v}_2)^{(4)},(\bar{u}_1)^{(4)},(\bar{u}_2)^{(4)}$ :<br><br>By $(\bar{v}_1)^{(4)}>0$ , $(\bar{v}_2)^{(4)}<0$ and respectively $(\bar{u}_1)^{(4)}>0$ , $(\bar{u}_2)^{(4)}<0$ the roots of the equations $(a_{25})^{(4)}\left(v^{(4)}\right)^2+(\sigma_2)^{(4)}v^{(4)}-(a_{24})^{(4)}=0$ and $(b_{25})^{(4)}\left(u^{(4)}\right)^2+(\tau_2)^{(4)}u^{(4)}-(b_{24})^{(4)}=0$<br>**Definition of** $(m_1)^{(4)}$ , $(m_2)^{(4)}$ , $(\mu_1)^{(4)},(\mu_2)^{(4)},(v_0)^{(4)}$ :-<br><br>If we define $(m_1)^{(4)}$ , $(m_2)^{(4)}$ , $(\mu_1)^{(4)},(\mu_2)^{(4)}$ by<br><br>$(m_2)^{(4)}=(v_0)^{(4)},(m_1)^{(4)}=(v_1)^{(4)}$, ***if*** $(v_0)^{(4)}<(v_1)^{(4)}$<br><br>$(m_2)^{(4)}=(v_1)^{(4)},(m_1)^{(4)}=(\bar{v}_1)^{(4)}$ , ***if*** $(v_4)^{(4)}<(v_0)^{(4)}<(\bar{v}_1)^{(4)}$,<br>and $\boxed{(v_0)^{(4)}=\frac{G_{24}^0}{G_{25}^0}}$<br><br>$(m_2)^{(4)}=(v_4)^{(4)},(m_1)^{(4)}=(v_0)^{(4)}$, ***if*** $(\bar{v}_4)^{(4)}<(v_0)^{(4)}$ | 330 |
| and analogously<br><br>$(\mu_2)^{(4)}=(u_0)^{(4)},(\mu_1)^{(4)}=(u_1)^{(4)}$, ***if*** $(u_0)^{(4)}<(u_1)^{(4)}$<br><br>$(\mu_2)^{(4)}=(u_1)^{(4)},(\mu_1)^{(4)}=(\bar{u}_1)^{(4)}$ , ***if*** $(u_1)^{(4)}<(u_0)^{(4)}<(\bar{u}_1)^{(4)}$,<br>and $\boxed{(u_0)^{(4)}=\frac{T_{24}^0}{T_{25}^0}}$<br><br>$(\mu_2)^{(4)}=(u_1)^{(4)},(\mu_1)^{(4)}=(u_0)^{(4)}$, ***if*** $(\bar{u}_1)^{(4)}<(u_0)^{(4)}$ where $(u_1)^{(4)},(\bar{u}_1)^{(4)}$ | 331 |
| Then the solution of global equations satisfies the inequalities<br><br>$G_{24}^0e^{((S_1)^{(4)}-(p_{24})^{(4)})t}\le G_{24}(t)\le G_{24}^0e^{(S_1)^{(4)}t}$<br><br>where $(p_i)^{(4)}$ is defined by equation | 332 |
| $\frac{1}{(m_1)^{(4)}}G_{24}^0e^{((S_1)^{(4)}-(p_{24})^{(4)})t}\le G_{25}(t)\le\frac{1}{(m_2)^{(4)}}G_{24}^0e^{(S_1)^{(4)}t}$ | 333 |
| $\left(\frac{(a_{26})^{(4)}G_{24}^0}{(m_1)^{(4)}((S_1)^{(4)}-(p_{24})^{(4)}-(S_2)^{(4)})}\left[e^{((S_1)^{(4)}-(p_{24})^{(4)})t}-e^{-(S_2)^{(4)}t}\right]+G_{26}^0e^{-(S_2)^{(4)}t}\le G_{26}(t)\le\right.$<br>*(a26)4G240(m2)4(S1)4−(a26′)4e(S1)4t−e−(a26′)4t+ G260e−(a26′)4t* | 334 |
| $\boxed{T_{24}^0e^{(R_1)^{(4)}t}\le T_{24}(t)\le T_{24}^0e^{((R_1)^{(4)}+(r_{24})^{(4)})t}}$ | |
| $\frac{1}{(\mu_1)^{(4)}}T_{24}^0e^{(R_1)^{(4)}t}\le T_{24}(t)\le\frac{1}{(\mu_2)^{(4)}}T_{24}^0e^{((R_1)^{(4)}+(r_{24})^{(4)})t}$ | 335 |
| $\frac{(b_{26})^{(4)}T_{24}^0}{(\mu_1)^{(4)}((R_1)^{(4)}-(b_{26}^{'})^{(4)})}\left[e^{(R_1)^{(4)}t}-e^{-(b_{26}^{'})^{(4)}t}\right]+T_{26}^0e^{-(b_{26}^{'})^{(4)}t}\le T_{26}(t)\le$<br><br>$\frac{(a_{26})^{(4)}T_{24}^0}{(\mu_2)^{(4)}((R_1)^{(4)}+(r_{24})^{(4)}+(R_2)^{(4)})}\left[e^{((R_1)^{(4)}+(r_{24})^{(4)})t}-e^{-(R_2)^{(4)}t}\right]+T_{26}^0e^{-(R_2)^{(4)}t}$ | 336 |

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| **Definition of** $(S_1)^{(4)}, (S_2)^{(4)}, (R_1)^{(4)}, (R_2)^{(4)}$:-<br>Where $(S_1)^{(4)} = (a_{24})^{(4)}(m_2)^{(4)} - (a'_{24})^{(4)}$<br>$(S_2)^{(4)} = (a_{26})^{(4)} - (p_{26})^{(4)}$<br>$(R_1)^{(4)} = (b_{24})^{(4)}(\mu_2)^{(4)} - (b'_{24})^{(4)}$<br>$(R_2)^{(4)} = (b'_{26})^{(4)} - (r_{26})^{(4)}$ | 337 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ :<br>$(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ four constants satisfying<br>$-(\sigma_2)^{(5)} \le -(a'_{28})^{(5)} + (a'_{29})^{(5)} - (a''_{28})^{(5)}(T_{29}, t) + (a''_{29})^{(5)}(T_{29}, t) \le -(\sigma_1)^{(5)}$<br>$-(\tau_2)^{(5)} \le -(b'_{28})^{(5)} + (b'_{29})^{(5)} - (b''_{28})^{(5)}((G_{31}), t) - (b''_{29})^{(5)}((G_{31}), t) \le -(\tau_1)^{(5)}$ | 338 |
| **Definition of** $(v_1)^{(5)}, (v_2)^{(5)}, (u_1)^{(5)}, (u_2)^{(5)}, v^{(5)}, u^{(5)}$ :<br>By $(v_1)^{(5)} > 0$, $(v_2)^{(5)} < 0$ and respectively $(u_1)^{(5)} > 0$, $(u_2)^{(5)} < 0$ the roots of the equations $(a_{29})^{(5)}\left(v^{(5)}\right)^2 + (\sigma_1)^{(5)}v^{(5)} - (a_{28})^{(5)} = 0$<br>and $(b_{29})^{(5)}\left(u^{(5)}\right)^2 + (\tau_1)^{(5)}u^{(5)} - (b_{28})^{(5)} = 0$ and | 339 |
| **Definition of** $(\bar{v}_1)^{(5)},, (\bar{v}_2)^{(5)}, (\bar{u}_1)^{(5)}, (\bar{u}_2)^{(5)}$ :<br>By $(\bar{v}_1)^{(5)} > 0$, $(\bar{v}_2)^{(5)} < 0$ and respectively $(\bar{u}_1)^{(5)} > 0$, $(\bar{u}_2)^{(5)} < 0$ the roots of the equations $(a_{29})^{(5)}\left(v^{(5)}\right)^2 + (\sigma_2)^{(5)}v^{(5)} - (a_{28})^{(5)} = 0$<br>and $(b_{29})^{(5)}\left(u^{(5)}\right)^2 + (\tau_2)^{(5)}u^{(5)} - (b_{28})^{(5)} = 0$<br>**Definition of** $(m_1)^{(5)}, (m_2)^{(5)}, (\mu_1)^{(5)}, (\mu_2)^{(5)}, (v_0)^{(5)}$ :-<br>If we define $(m_1)^{(5)}, (m_2)^{(5)}, (\mu_1)^{(5)}, (\mu_2)^{(5)}$ by<br>$(m_2)^{(5)} = (v_0)^{(5)}, (m_1)^{(5)} = (v_1)^{(5)}$, ***if*** $(v_0)^{(5)} < (v_1)^{(5)}$<br>$(m_2)^{(5)} = (v_1)^{(5)}, (m_1)^{(5)} = (\bar{v}_1)^{(5)}$, ***if*** $(v_1)^{(5)} < (v_0)^{(5)} < (\bar{v}_1)^{(5)}$,<br>and $\boxed{(v_0)^{(5)} = \frac{G^0_{28}}{G^0_{29}}}$<br>$(m_2)^{(5)} = (v_1)^{(5)}, (m_1)^{(5)} = (v_0)^{(5)}$, ***if*** $(\bar{v}_1)^{(5)} < (v_0)^{(5)}$ | 340 |
| and analogously<br>$(\mu_2)^{(5)} = (u_0)^{(5)}, (\mu_1)^{(5)} = (u_1)^{(5)}$, ***if*** $(u_0)^{(5)} < (u_1)^{(5)}$ | 341 |

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| $(\mu_2)^{(5)} = (u_1)^{(5)}, (\mu_1)^{(5)} = (\bar{u}_1)^{(5)}$ , $\boldsymbol{if}$ $(u_1)^{(5)} < (u_0)^{(5)} < (\bar{u}_1)^{(5)}$, and $\boxed{(u_0)^{(5)} = \frac{T_{28}^0}{T_{29}^0}}$ $(\mu_2)^{(5)} = (u_1)^{(5)}, (\mu_1)^{(5)} = (u_0)^{(5)}, \boldsymbol{if}$ $(\bar{u}_1)^{(5)} < (u_0)^{(5)}$ where $(u_1)^{(5)}, (\bar{u}_1)^{(5)}$ | |
| Then the solution of global equations satisfies the inequalities $G_{28}^0 e^{((S_1)^{(5)} - (p_{28})^{(5)})t} \le G_{28}(t) \le G_{28}^0 e^{(S_1)^{(5)} t}$ where $(p_i)^{(5)}$ is defined by equation | 342 |
| $\frac{1}{(m_5)^{(5)}} G_{28}^0 e^{((S_1)^{(5)} - (p_{28})^{(5)})t} \le G_{29}(t) \le \frac{1}{(m_2)^{(5)}} G_{28}^0 e^{(S_1)^{(5)} t}$ | 343 |
| $\left(\frac{(a_{30})^{(5)} G_{28}^0}{(m_1)^{(5)}((S_1)^{(5)} - (p_{28})^{(5)} - (S_2)^{(5)})}\left[e^{((S_1)^{(5)} - (p_{28})^{(5)})t} - e^{-(S_2)^{(5)} t}\right] + G_{30}^0 e^{-(S_2)^{(5)} t} \le G_{30}(t) \le$ *(a30)5G280(m2)5(S1)5−(a30′)5e(S1)5t−e−(a30′)5t+ G300e−(a30′)5t* | 344 |
| $\boxed{T_{28}^0 e^{(R_1)^{(5)} t} \le T_{28}(t) \le T_{28}^0 e^{((R_1)^{(5)} + (r_{28})^{(5)})t}}$ | 345 |
| $\frac{1}{(\mu_1)^{(5)}} T_{28}^0 e^{(R_1)^{(5)} t} \le T_{28}(t) \le \frac{1}{(\mu_2)^{(5)}} T_{28}^0 e^{((R_1)^{(5)} + (r_{28})^{(5)})t}$ | 346 |
| $\frac{(b_{30})^{(5)} T_{28}^0}{(\mu_1)^{(5)}((R_1)^{(5)} - (b_{30}^{'})^{(5)})}\left[e^{(R_1)^{(5)} t} - e^{-(b_{30}^{'})^{(5)} t}\right] + T_{30}^0 e^{-(b_{30}^{'})^{(5)} t} \le T_{30}(t) \le$ $\frac{(a_{30})^{(5)} T_{28}^0}{(\mu_2)^{(5)}((R_1)^{(5)} + (r_{28})^{(5)} + (R_2)^{(5)})}\left[e^{((R_1)^{(5)} + (r_{28})^{(5)})t} - e^{-(R_2)^{(5)} t}\right] + T_{30}^0 e^{-(R_2)^{(5)} t}$ | 347 |
| **Definition of** $(S_1)^{(5)}, (S_2)^{(5)}, (R_1)^{(5)}, (R_2)^{(5)}$:- Where $(S_1)^{(5)} = (a_{28})^{(5)}(m_2)^{(5)} - (a_{28}^{'})^{(5)}$ $(S_2)^{(5)} = (a_{30})^{(5)} - (p_{30})^{(5)}$ $(R_1)^{(5)} = (b_{28})^{(5)}(\mu_2)^{(5)} - (b_{28}^{'})^{(5)}$ $(R_2)^{(5)} = (b_{30}^{'})^{(5)} - (r_{30})^{(5)}$ | 348 |
| **Behavior of the solutions of equation** **Theorem 2:** If we denote and define **Definition of** $(\sigma_1)^{(6)}, (\sigma_2)^{(6)}, (\tau_1)^{(6)}, (\tau_2)^{(6)}$ : $(\sigma_1)^{(6)}, (\sigma_2)^{(6)}, (\tau_1)^{(6)}, (\tau_2)^{(6)}$ four constants satisfying $-(\sigma_2)^{(6)} \le -(a_{32}^{'})^{(6)} + (a_{33}^{'})^{(6)} - (a_{32}^{''})^{(6)}(T_{33}, t) + (a_{33}^{''})^{(6)}(T_{33}, t) \le -(\sigma_1)^{(6)}$ | 349 |

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| $-(\tau_2)^{(6)} \le -(b'_{32})^{(6)} + (b'_{33})^{(6)} - (b''_{32})^{(6)}\big((G_{35}),t\big) - (b''_{33})^{(6)}\big((G_{35}),t\big) \le -(\tau_1)^{(6)}$ | |
| **Definition of** $(v_1)^{(6)},(v_2)^{(6)},(u_1)^{(6)},(u_2)^{(6)},v^{(6)},u^{(6)}$ :<br><br>By $(v_1)^{(6)} > 0$ , $(v_2)^{(6)} < 0$ and respectively $(u_1)^{(6)} > 0$ , $(u_2)^{(6)} < 0$ the roots of the equations $(a_{33})^{(6)}\big(v^{(6)}\big)^2 + (\sigma_1)^{(6)}v^{(6)} - (a_{32})^{(6)} = 0$<br>and $(b_{33})^{(6)}\big(u^{(6)}\big)^2 + (\tau_1)^{(6)}u^{(6)} - (b_{32})^{(6)} = 0$ and | 350 |
| **Definition of** $(\bar{v}_1)^{(6)},,(\bar{v}_2)^{(6)},(\bar{u}_1)^{(6)},(\bar{u}_2)^{(6)}$ :<br><br>By $(\bar{v}_1)^{(6)} > 0$ , $(\bar{v}_2)^{(6)} < 0$ and respectively $(\bar{u}_1)^{(6)} > 0$ , $(\bar{u}_2)^{(6)} < 0$ the roots of the equations $(a_{33})^{(6)}\big(v^{(6)}\big)^2 + (\sigma_2)^{(6)}v^{(6)} - (a_{32})^{(6)} = 0$<br>and $(b_{33})^{(6)}\big(u^{(6)}\big)^2 + (\tau_2)^{(6)}u^{(6)} - (b_{32})^{(6)} = 0$<br>**Definition of** $(m_1)^{(6)}$ , $(m_2)^{(6)}$ , $(\mu_1)^{(6)},(\mu_2)^{(6)},(v_0)^{(6)}$ :-<br><br>If we define $(m_1)^{(6)}$ , $(m_2)^{(6)}$ , $(\mu_1)^{(6)},(\mu_2)^{(6)}$ by<br><br>$(m_2)^{(6)} = (v_0)^{(6)}, (m_1)^{(6)} = (v_1)^{(6)}$, ***if*** $(v_0)^{(6)} < (v_1)^{(6)}$<br><br>$(m_2)^{(6)} = (v_1)^{(6)}, (m_1)^{(6)} = (\bar{v}_6)^{(6)}$ , ***if*** $(v_1)^{(6)} < (v_0)^{(6)} < (\bar{v}_1)^{(6)}$,<br>and $\boxed{(v_0)^{(6)} = \frac{G^0_{32}}{G^0_{33}}}$<br><br>$(m_2)^{(6)} = (v_1)^{(6)}, (m_1)^{(6)} = (v_0)^{(6)}$, ***if*** $(\bar{v}_1)^{(6)} < (v_0)^{(6)}$ | 351 |
| and analogously<br><br>$(\mu_2)^{(6)} = (u_0)^{(6)}, (\mu_1)^{(6)} = (u_1)^{(6)}$, ***if*** $(u_0)^{(6)} < (u_1)^{(6)}$<br><br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (\bar{u}_1)^{(6)}$ , ***if*** $(u_1)^{(6)} < (u_0)^{(6)} < (\bar{u}_1)^{(6)}$,<br>and $\boxed{(u_0)^{(6)} = \frac{T^0_{32}}{T^0_{33}}}$<br><br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (u_0)^{(6)}$, ***if*** $(\bar{u}_1)^{(6)} < (u_0)^{(6)}$ where $(u_1)^{(6)}, (\bar{u}_1)^{(6)}$ | 352 |
| Then the solution of global equations satisfies the inequalities<br><br>$G^0_{32}e^{((S_1)^{(6)}-(p_{32})^{(6)})t} \le G_{32}(t) \le G^0_{32}e^{(S_1)^{(6)}t}$<br><br>where $(p_i)^{(6)}$ is defined by equation | 353 |
| $\frac{1}{(m_1)^{(6)}}G^0_{32}e^{((S_1)^{(6)}-(p_{32})^{(6)})t} \le G_{33}(t) \le \frac{1}{(m_2)^{(6)}}G^0_{32}e^{(S_1)^{(6)}t}$ | 354 |
| $\left(\frac{(a_{34})^{(6)}G^0_{32}}{(m_1)^{(6)}((S_1)^{(6)}-(p_{32})^{(6)}-(S_2)^{(6)})}\left[e^{((S_1)^{(6)}-(p_{32})^{(6)})t} - e^{-(S_2)^{(6)}t}\right] + G^0_{34}e^{-(S_2)^{(6)}t} \le G_{34}(t) \le$<br>*(a34)6G320(m2)6(S1)6−(a34′)6e(S1)6t−e−(a34′)6t+ G340e−(a34′)6t* | 355 |

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| $T_{32}^{0}e^{(R_1)^{(6)}t} \le T_{32}(t) \le T_{32}^{0}e^{((R_1)^{(6)}+(r_{32})^{(6)})t}$ | 356 |
| $\frac{1}{(\mu_1)^{(6)}}T_{32}^{0}e^{(R_1)^{(6)}t} \le T_{32}(t) \le \frac{1}{(\mu_2)^{(6)}}T_{32}^{0}e^{((R_1)^{(6)}+(r_{32})^{(6)})t}$ | 357 |
| $\frac{(b_{34})^{(6)}T_{32}^{0}}{(\mu_1)^{(6)}((R_1)^{(6)}-(b_{34}^{'})^{(6)})}\left[e^{(R_1)^{(6)}t}-e^{-(b_{34}^{'})^{(6)}t}\right]+T_{34}^{0}e^{-(b_{34}^{'})^{(6)}t} \le T_{34}(t) \le$ $\frac{(a_{34})^{(6)}T_{32}^{0}}{(\mu_2)^{(6)}((R_1)^{(6)}+(r_{32})^{(6)}+(R_2)^{(6)})}\left[e^{((R_1)^{(6)}+(r_{32})^{(6)})t}-e^{-(R_2)^{(6)}t}\right]+T_{34}^{0}e^{-(R_2)^{(6)}t}$ | 358 |
| **Definition of** $(S_1)^{(6)},(S_2)^{(6)},(R_1)^{(6)},(R_2)^{(6)}$:- Where $(S_1)^{(6)}=(a_{32})^{(6)}(m_2)^{(6)}-(a_{32}^{'})^{(6)}$ $(S_2)^{(6)}=(a_{34})^{(6)}-(p_{34})^{(6)}$ $(R_1)^{(6)}=(b_{32})^{(6)}(\mu_2)^{(6)}-(b_{32}^{'})^{(6)}$ $(R_2)^{(6)}=(b_{34}^{'})^{(6)}-(r_{34})^{(6)}$ | 359 |
| **Behavior of the solutions of equation** **Theorem 2:** If we denote and define **Definition of** $(\sigma_1)^{(7)},(\sigma_2)^{(7)},(\tau_1)^{(7)},(\tau_2)^{(7)}$: $(\sigma_1)^{(7)},(\sigma_2)^{(7)},(\tau_1)^{(7)},(\tau_2)^{(7)}$ four constants satisfying $-(\sigma_2)^{(7)} \le -(a_{36}^{'})^{(7)}+(a_{37}^{'})^{(7)}-(a_{36}^{''})^{(7)}(T_{37},t)+(a_{37}^{''})^{(7)}(T_{37},t) \le -(\sigma_1)^{(7)}$ $-(\tau_2)^{(7)} \le -(b_{36}^{'})^{(7)}+(b_{37}^{'})^{(7)}-(b_{36}^{''})^{(7)}((G_{39}),t)-(b_{37}^{''})^{(7)}((G_{39}),t) \le -(\tau_1)^{(7)}$ | |
| **Definition of** $(v_1)^{(7)},(v_2)^{(7)},(u_1)^{(7)},(u_2)^{(7)},v^{(7)},u^{(7)}$: By $(v_1)^{(7)}>0,(v_2)^{(7)}<0$ and respectively $(u_1)^{(7)}>0,(u_2)^{(7)}<0$ the roots of the equations $(a_{37})^{(7)}\left(v^{(7)}\right)^2+(\sigma_1)^{(7)}v^{(7)}-(a_{36})^{(7)}=0$ and $(b_{37})^{(7)}\left(u^{(7)}\right)^2+(\tau_1)^{(7)}u^{(7)}-(b_{36})^{(7)}=0$ and | 361 |
| **Definition of** $(\bar{v}_1)^{(7)},,(\bar{v}_2)^{(7)},(\bar{u}_1)^{(7)},(\bar{u}_2)^{(7)}$: By $(\bar{v}_1)^{(7)}>0,(\bar{v}_2)^{(7)}<0$ and respectively $(\bar{u}_1)^{(7)}>0,(\bar{u}_2)^{(7)}<0$ the roots of the equations $(a_{37})^{(7)}\left(v^{(7)}\right)^2+(\sigma_2)^{(7)}v^{(7)}-(a_{36})^{(7)}=0$ and $(b_{37})^{(7)}\left(u^{(7)}\right)^2+(\tau_2)^{(7)}u^{(7)}-(b_{36})^{(7)}=0$ **Definition of** $(m_1)^{(7)},(m_2)^{(7)},(\mu_1)^{(7)},(\mu_2)^{(7)},(v_0)^{(7)}$:- | 362 |

If we define $(m_1)^{(7)}, (m_2)^{(7)}, (\mu_1)^{(7)}, (\mu_2)^{(7)}$ by

$$(m_2)^{(7)} = (\nu_0)^{(7)}, (m_1)^{(7)} = (\nu_1)^{(7)}, \ \boldsymbol{if}\ (\nu_0)^{(7)} < (\nu_1)^{(7)}$$

$$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\bar{\nu}_1)^{(7)}, \boldsymbol{if}\ (\nu_1)^{(7)} < (\nu_0)^{(7)} < (\bar{\nu}_1)^{(7)},$$

and $\boxed{(\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0}}$

$$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\nu_0)^{(7)}, \ \boldsymbol{if}\ (\bar{\nu}_1)^{(7)} < (\nu_0)^{(7)}$$

and analogously

$$(\mu_2)^{(7)} = (u_0)^{(7)}, (\mu_1)^{(7)} = (u_1)^{(7)}, \ \boldsymbol{if}\ (u_0)^{(7)} < (u_1)^{(7)}$$

$$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (\bar{u}_1)^{(7)}, \boldsymbol{if}\ (u_1)^{(7)} < (u_0)^{(7)} < (\bar{u}_1)^{(7)},$$

and $\boxed{(u_0)^{(7)} = \frac{T_{36}^0}{T_{37}^0}}$

$$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (u_0)^{(7)}, \boldsymbol{if}\ (\bar{u}_1)^{(7)} < (u_0)^{(7)}$$ where $(u_1)^{(7)}, (\bar{u}_1)^{(7)}$

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Then the solution of global equations satisfies the inequalities

$$G_{36}^0 e^{((S_1)^{(7)} - (p_{36})^{(7)})t} \le G_{36}(t) \le G_{36}^0 e^{(S_1)^{(7)} t}$$

where $(p_i)^{(7)}$ is defined by equation

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$$\frac{1}{(m_7)^{(7)}} G_{36}^0 e^{((S_1)^{(7)} - (p_{36})^{(7)})t} \le G_{37}(t) \le \frac{1}{(m_2)^{(7)}} G_{36}^0 e^{(S_1)^{(7)} t}$$

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$$\left(\frac{(a_{38})^{(7)} G_{36}^0}{(m_1)^{(7)}((S_1)^{(7)} - (p_{36})^{(7)} - (S_2)^{(7)})}\left[e^{((S_1)^{(7)} - (p_{36})^{(7)})t} - e^{-(S_2)^{(7)} t}\right] + G_{38}^0 e^{-(S_2)^{(7)} t} \le G_{38}(t) \le \right.$$

$$\left. \frac{(a_{38})^{(7)} G_{36}^0}{(m_2)^{(7)}((S_1)^{(7)} - (a_{38}')^{(7)})}\left[e^{(S_1)^{(7)} t} - e^{-(a_{38}')^{(7)} t}\right] + G_{38}^0 e^{-(a_{38}')^{(7)} t}\right)$$

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$$\boxed{T_{36}^0 e^{(R_1)^{(7)} t} \le T_{36}(t) \le T_{36}^0 e^{((R_1)^{(7)} + (r_{36})^{(7)})t}}$$

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$$\frac{1}{(\mu_1)^{(7)}} T_{36}^0 e^{(R_1)^{(7)} t} \le T_{36}(t) \le \frac{1}{(\mu_2)^{(7)}} T_{36}^0 e^{((R_1)^{(7)} + (r_{36})^{(7)})t}$$

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$$\frac{(b_{38})^{(7)} T_{36}^0}{(\mu_1)^{(7)}((R_1)^{(7)} - (b_{38}')^{(7)})}\left[e^{(R_1)^{(7)} t} - e^{-(b_{38}')^{(7)} t}\right] + T_{38}^0 e^{-(b_{38}')^{(7)} t} \le T_{38}(t) \le$$

$$\frac{(a_{38})^{(7)} T_{36}^0}{(\mu_2)^{(7)}((R_1)^{(7)} + (r_{36})^{(7)} + (R_2)^{(7)})}\left[e^{((R_1)^{(7)} + (r_{36})^{(7)})t} - e^{-(R_2)^{(7)} t}\right] + T_{38}^0 e^{-(R_2)^{(7)} t}$$

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| **Definition of** $(S_1)^{(7)}, (S_2)^{(7)}, (R_1)^{(7)}, (R_2)^{(7)}$:-<br><br>Where $(S_1)^{(7)} = (a_{36})^{(7)}(m_2)^{(7)} - (a'_{36})^{(7)}$<br><br>$(S_2)^{(7)} = (a_{38})^{(7)} - (p_{38})^{(7)}$<br><br>$(R_1)^{(7)} = (b_{36})^{(7)}(\mu_2)^{(7)} - (b'_{36})^{(7)}$<br><br>$(R_2)^{(7)} = (b'_{38})^{(7)} - (r_{38})^{(7)}$ | 370 |
| **Behavior of the solutions of equation**<br><br>**Theorem 2:** If we denote and define<br><br>**Definition of** $(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ :<br><br>$(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ four constants satisfying<br><br>$-(\sigma_2)^{(8)} \le -(a'_{40})^{(8)} + (a'_{41})^{(8)} - (a''_{40})^{(8)}(T_{41}, t) + (a''_{41})^{(8)}(T_{41}, t) \le -(\sigma_1)^{(8)}$<br><br>$-(\tau_2)^{(8)} \le -(b'_{40})^{(8)} + (b'_{41})^{(8)} - (b''_{40})^{(8)}\big((G_{43}), t\big) - (b''_{41})^{(8)}\big((G_{43}), t\big) \le -(\tau_1)^{(8)}$ | 371 |
| **Definition of** $(v_1)^{(8)}, (v_2)^{(8)}, (u_1)^{(8)}, (u_2)^{(8)}, v^{(8)}, u^{(8)}$ :<br><br>By $(v_1)^{(8)} > 0, (v_2)^{(8)} < 0$ and respectively $(u_1)^{(8)} > 0, (u_2)^{(8)} < 0$ the roots of the equations $(a_{41})^{(8)}\big(v^{(8)}\big)^2 + (\sigma_1)^{(8)}v^{(8)} - (a_{40})^{(8)} = 0$<br>and $(b_{41})^{(8)}\big(u^{(8)}\big)^2 + (\tau_1)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$ and | 372 |
| **Definition of** $(\bar{v}_1)^{(8)}, , (\bar{v}_2)^{(8)}, (\bar{u}_1)^{(8)}, (\bar{u}_2)^{(8)}$ :<br><br>By $(\bar{v}_1)^{(8)} > 0, (\bar{v}_2)^{(8)} < 0$ and respectively $(\bar{u}_1)^{(8)} > 0, (\bar{u}_2)^{(8)} < 0$ the<br><br>roots of the equations $(a_{41})^{(8)}\big(v^{(8)}\big)^2 + (\sigma_2)^{(8)}v^{(8)} - (a_{40})^{(8)} = 0$<br><br>and $(b_{41})^{(8)}\big(u^{(8)}\big)^2 + (\tau_2)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$<br><br>**Definition of** $(m_1)^{(8)}, (m_2)^{(8)}, (\mu_1)^{(8)}, (\mu_2)^{(8)}, (v_0)^{(8)}$ :-<br><br>If we define $(m_1)^{(8)}, (m_2)^{(8)}, (\mu_1)^{(8)}, (\mu_2)^{(8)}$ by<br><br>$(m_2)^{(8)} = (v_0)^{(8)}, (m_1)^{(8)} = (v_1)^{(8)}$, ***if*** $(v_0)^{(8)} < (v_1)^{(8)}$<br><br>$(m_2)^{(8)} = (v_1)^{(8)}, (m_1)^{(8)} = (\bar{v}_1)^{(8)}$, ***if*** $(v_1)^{(8)} < (v_0)^{(8)} < (\bar{v}_1)^{(8)}$, | |

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| and $\boxed{(v_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}}$<br><br>$(m_2)^{(8)} = (v_1)^{(8)}, (m_1)^{(8)} = (v_0)^{(8)},\ \boldsymbol{if}\ (\bar{v}_1)^{(8)} < (v_0)^{(8)}$ | |
| and analogously<br><br>$(\mu_2)^{(8)} = (u_0)^{(8)}, (\mu_1)^{(8)} = (u_1)^{(8)},\ \boldsymbol{if}\ (u_0)^{(8)} < (u_1)^{(8)}$<br><br>$(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (\bar{u}_1)^{(8)}, \boldsymbol{if}\ (u_1)^{(8)} < (u_0)^{(8)} < (\bar{u}_1)^{(8)},$<br>and $\boxed{(u_0)^{(8)} = \frac{T_{40}^0}{T_{41}^0}}$<br><br>$(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (u_0)^{(8)}, \boldsymbol{if}\ (\bar{u}_1)^{(8)} < (u_0)^{(8)}$ where $(u_1)^{(8)}, (\bar{u}_1)^{(8)}$ | 374 |
| Then the solution of global equations satisfies the inequalities<br><br>$G_{40}^0 e^{((S_1)^{(8)} - (p_{40})^{(8)})t} \le G_{40}(t) \le G_{40}^0 e^{(S_1)^{(8)} t}$<br><br>where $(p_i)^{(8)}$ is defined by equation | 375 |
| $\frac{1}{(m_1)^{(8)}} G_{40}^0 e^{((S_1)^{(8)} - (p_{40})^{(8)})t} \le G_{41}(t) \le \frac{1}{(m_2)^{(8)}} G_{40}^0 e^{(S_1)^{(8)} t}$ | 376 |
| $\left(\frac{(a_{42})^{(8)} G_{40}^0}{(m_1)^{(8)}((S_1)^{(8)} - (p_{40})^{(8)} - (S_2)^{(8)})}\left[e^{((S_1)^{(8)} - (p_{40})^{(8)})t} - e^{-(S_2)^{(8)} t}\right] + G_{42}^0 e^{-(S_2)^{(8)} t} \le G_{42}(t) \le \right.$<br>$\left.\frac{(a_{42})^{(8)} G_{40}^0}{(m_2)^{(8)}((S_1)^{(8)} - (a'_{42})^{(8)})}\left[e^{(S_1)^{(8)} t} - e^{-(a'_{42})^{(8)} t}\right] + G_{42}^0 e^{-(a'_{42})^{(8)} t}\right)$ | 377 |
| $\boxed{T_{40}^0 e^{(R_1)^{(8)} t} \le T_{40}(t) \le T_{40}^0 e^{((R_1)^{(8)} + (r_{40})^{(8)})t}}$ | 378 |
| $\frac{1}{(\mu_1)^{(8)}} T_{40}^0 e^{(R_1)^{(8)} t} \le T_{40}(t) \le \frac{1}{(\mu_2)^{(8)}} T_{40}^0 e^{((R_1)^{(8)} + (r_{40})^{(8)})t}$ | 379 |
| $\frac{(b_{42})^{(8)} T_{40}^0}{(\mu_1)^{(8)}((R_1)^{(8)} - (b'_{42})^{(8)})}\left[e^{(R_1)^{(8)} t} - e^{-(b'_{42})^{(8)} t}\right] + T_{42}^0 e^{-(b'_{42})^{(8)} t} \le T_{42}(t) \le$<br><br>$\frac{(a_{42})^{(8)} T_{40}^0}{(\mu_2)^{(8)}((R_1)^{(8)} + (r_{40})^{(8)} + (R_2)^{(8)})}\left[e^{((R_1)^{(8)} + (r_{40})^{(8)})t} - e^{-(R_2)^{(8)} t}\right] + T_{42}^0 e^{-(R_2)^{(8)} t}$ | 380 |
| **Definition of** $(S_1)^{(8)}, (S_2)^{(8)}, (R_1)^{(8)}, (R_2)^{(8)}$:-<br><br>Where $(S_1)^{(8)} = (a_{40})^{(8)}(m_2)^{(8)} - (a'_{40})^{(8)}$<br><br>$(S_2)^{(8)} = (a_{42})^{(8)} - (p_{42})^{(8)}$ | 381 |

$(R_1)^{(8)} = (b_{40})^{(8)}(\mu_2)^{(8)} - (b'_{40})^{(8)}$

$(R_2)^{(8)} = (b'_{42})^{(8)} - (r_{42})^{(8)}$

**Behavior of the solutions of equation 37 to 92** 382

**Theorem 2:** If we denote and define

**Definition of** $(\sigma_1)^{(9)}, (\sigma_2)^{(9)}, (\tau_1)^{(9)}, (\tau_2)^{(9)}$ :

$(\sigma_1)^{(9)}, (\sigma_2)^{(9)}, (\tau_1)^{(9)}, (\tau_2)^{(9)}$ four constants satisfying

$-(\sigma_2)^{(9)} \le -(a'_{44})^{(9)} + (a'_{45})^{(9)} - (a''_{44})^{(9)}(T_{45}, t) + (a''_{45})^{(9)}(T_{45}, t) \le -(\sigma_1)^{(9)}$

$-(\tau_2)^{(9)} \le -(b'_{44})^{(9)} + (b'_{45})^{(9)} - (b''_{44})^{(9)}\big((G_{47}), t\big) - (b''_{45})^{(9)}\big((G_{47}), t\big) \le -(\tau_1)^{(9)}$

**Definition of** $(\nu_1)^{(9)}, (\nu_2)^{(9)}, (u_1)^{(9)}, (u_2)^{(9)}, \nu^{(9)}, u^{(9)}$ :

By $(\nu_1)^{(9)} > 0$ , $(\nu_2)^{(9)} < 0$ and respectively $(u_1)^{(9)} > 0$ , $(u_2)^{(9)} < 0$ the roots of the equations $(a_{45})^{(9)}\big(\nu^{(9)}\big)^2 + (\sigma_1)^{(9)}\nu^{(9)} - (a_{44})^{(9)} = 0$ and $(b_{45})^{(9)}\big(u^{(9)}\big)^2 + (\tau_1)^{(9)}u^{(9)} - (b_{44})^{(9)} = 0$ and

**Definition of** $(\bar{\nu}_1)^{(9)}, , (\bar{\nu}_2)^{(9)}, (\bar{u}_1)^{(9)}, (\bar{u}_2)^{(9)}$ :

By $(\bar{\nu}_1)^{(9)} > 0$ , $(\bar{\nu}_2)^{(9)} < 0$ and respectively $(\bar{u}_1)^{(9)} > 0$ , $(\bar{u}_2)^{(9)} < 0$ the roots of the equations $(a_{45})^{(9)}\big(\nu^{(9)}\big)^2 + (\sigma_2)^{(9)}\nu^{(9)} - (a_{44})^{(9)} = 0$ and $(b_{45})^{(9)}\big(u^{(9)}\big)^2 + (\tau_2)^{(9)}u^{(9)} - (b_{44})^{(9)} = 0$

**Definition of** $(m_1)^{(9)}, (m_2)^{(9)}, (\mu_1)^{(9)}, (\mu_2)^{(9)}, (\nu_0)^{(9)}$ :-

If we define $(m_1)^{(9)}, (m_2)^{(9)}, (\mu_1)^{(9)}, (\mu_2)^{(9)}$ by

$(m_2)^{(9)} = (\nu_0)^{(9)}, (m_1)^{(9)} = (\nu_1)^{(9)}$, ***if*** $(\nu_0)^{(9)} < (\nu_1)^{(9)}$

$(m_2)^{(9)} = (\nu_1)^{(9)}, (m_1)^{(9)} = (\bar{\nu}_1)^{(9)}$ , ***if*** $(\nu_1)^{(9)} < (\nu_0)^{(9)} < (\bar{\nu}_1)^{(9)}$,

and $\boxed{(\nu_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}}}$

$(m_2)^{(9)} = (\nu_1)^{(9)}, (m_1)^{(9)} = (\nu_0)^{(9)}$, ***if*** $(\bar{\nu}_1)^{(9)} < (\nu_0)^{(9)}$

and analogously

$(\mu_2)^{(9)} = (u_0)^{(9)}, (\mu_1)^{(9)} = (u_1)^{(9)}$, ***if*** $(u_0)^{(9)} < (u_1)^{(9)}$

$(\mu_2)^{(9)} = (u_1)^{(9)}, (\mu_1)^{(9)} = (\bar{u}_1)^{(9)}$ , ***if*** $(u_1)^{(9)} < (u_0)^{(9)} < (\bar{u}_1)^{(9)}$,

and $\boxed{(u_0)^{(9)} = \frac{T^0_{44}}{T^0_{45}}}$

$(\mu_2)^{(9)} = (u_1)^{(9)}, (\mu_1)^{(9)} = (u_0)^{(9)}$, ***if*** $(\bar{u}_1)^{(9)} < (u_0)^{(9)}$ where $(u_1)^{(9)}, (\bar{u}_1)^{(9)}$ are defined by 59 and 69 respectively

Then the solution of 19,20,21,22,23 and 24 satisfies the inequalities

$$G_{44}^{0}e^{((S_1)^{(9)}-(p_{44})^{(9)})t} \le G_{44}(t) \le G_{44}^{0}e^{(S_1)^{(9)}t}$$

where $(p_i)^{(9)}$ is defined by equation 45

$$\frac{1}{(m_9)^{(9)}}G_{44}^{0}e^{((S_1)^{(9)}-(p_{44})^{(9)})t} \le G_{45}(t) \le \frac{1}{(m_2)^{(9)}}G_{44}^{0}e^{(S_1)^{(9)}t}$$

$$\left(\frac{(a_{46})^{(9)}G_{44}^{0}}{(m_1)^{(9)}((S_1)^{(9)}-(p_{44})^{(9)}-(S_2)^{(9)})}\left[e^{((S_1)^{(9)}-(p_{44})^{(9)})t}-e^{-(S_2)^{(9)}t}\right]+G_{46}^{0}e^{-(S_2)^{(9)}t} \le G_{46}(t) \le \right.$$

$$\left.\frac{(a_{46})^{(9)}G_{44}^{0}}{(m_2)^{(9)}((S_1)^{(9)}-(a_{46}^{'})^{(9)})}[e^{(S_1)^{(9)}t}-e^{-(a_{46}^{'})^{(9)}t}]+\ G_{46}^{0}e^{-(a_{46}^{'})^{(9)}t}\right)$$

$$\boxed{T_{44}^{0}e^{(R_1)^{(9)}t} \le T_{44}(t) \le T_{44}^{0}e^{((R_1)^{(9)}+(r_{44})^{(9)})t}}$$

$$\frac{1}{(\mu_1)^{(9)}}T_{44}^{0}e^{(R_1)^{(9)}t} \le T_{44}(t) \le \frac{1}{(\mu_2)^{(9)}}T_{44}^{0}e^{((R_1)^{(9)}+(r_{44})^{(9)})t}$$

$$\frac{(b_{46})^{(9)}T_{44}^{0}}{(\mu_1)^{(9)}((R_1)^{(9)}-(b_{46}^{'})^{(9)})}\left[e^{(R_1)^{(9)}t}-e^{-(b_{46}^{'})^{(9)}t}\right]+T_{46}^{0}e^{-(b_{46}^{'})^{(9)}t} \le T_{46}(t) \le$$

$$\frac{(a_{46})^{(9)}T_{44}^{0}}{(\mu_2)^{(9)}((R_1)^{(9)}+(r_{44})^{(9)}+(R_2)^{(9)})}\left[e^{((R_1)^{(9)}+(r_{44})^{(9)})t}-e^{-(R_2)^{(9)}t}\right]+T_{46}^{0}e^{-(R_2)^{(9)}t}$$

**Definition of** $(S_1)^{(9)}, (S_2)^{(9)}, (R_1)^{(9)}, (R_2)^{(9)}$:-

Where $(S_1)^{(9)} = (a_{44})^{(9)}(m_2)^{(9)} - (a_{44}^{'})^{(9)}$

$(S_2)^{(9)} = (a_{46})^{(9)} - (p_{46})^{(9)}$

$(R_1)^{(9)} = (b_{44})^{(9)}(\mu_2)^{(9)} - (b_{44}^{'})^{(9)}$

$(R_2)^{(9)} = (b_{46}^{'})^{(9)} - (r_{46})^{(9)}$

**Proof :** From global equations we obtain

$$\frac{dv^{(1)}}{dt} = (a_{13})^{(1)} - \left((a_{13}^{'})^{(1)} - (a_{14}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14},t)\right) - (a_{14}^{''})^{(1)}(T_{14},t)v^{(1)} - (a_{14})^{(1)}v^{(1)}$$

**Definition of** $v^{(1)}$ :- $\boxed{v^{(1)} = \frac{G_{13}}{G_{14}}}$

It follows

$$-\left((a_{14})^{(1)}(v^{(1)})^{2} + (\sigma_2)^{(1)}v^{(1)} - (a_{13})^{(1)}\right) \le \frac{dv^{(1)}}{dt} \le -\left((a_{14})^{(1)}(v^{(1)})^{2} + (\sigma_1)^{(1)}v^{(1)} - (a_{13})^{(1)}\right)$$

From which one obtains

**Definition of** $(\bar{v}_1)^{(1)}, (v_0)^{(1)}$ :-

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For $0 < \boxed{(\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}} < (\nu_1)^{(1)} < (\bar{\nu}_1)^{(1)}$

$$\nu^{(1)}(t) \geq \frac{(\nu_1)^{(1)} + (C)^{(1)}(\nu_2)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)} - (\nu_0)^{(1)}\right)t\right]}}{1 + (C)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)} - (\nu_0)^{(1)}\right)t\right]}} \quad , \quad \boxed{(C)^{(1)} = \frac{(\nu_1)^{(1)} - (\nu_0)^{(1)}}{(\nu_0)^{(1)} - (\nu_2)^{(1)}}}$$

it follows $(\nu_0)^{(1)} \leq \nu^{(1)}(t) \leq (\nu_1)^{(1)}$

In the same manner , we get

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$$\nu^{(1)}(t) \leq \frac{(\bar{\nu}_1)^{(1)} + (\bar{C})^{(1)}(\bar{\nu}_2)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}}{1 + (\bar{C})^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(1)} = \frac{(\bar{\nu}_1)^{(1)} - (\nu_0)^{(1)}}{(\nu_0)^{(1)} - (\bar{\nu}_2)^{(1)}}}$$

From which we deduce $(\nu_0)^{(1)} \leq \nu^{(1)}(t) \leq (\bar{\nu}_1)^{(1)}$

If $0 < (\nu_1)^{(1)} < (\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0} < (\bar{\nu}_1)^{(1)}$ we find like in the previous case,

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$$(\nu_1)^{(1)} \leq \frac{(\nu_1)^{(1)} + (C)^{(1)}(\nu_2)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)} - (\nu_2)^{(1)}\right)t\right]}}{1 + (C)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)} - (\nu_2)^{(1)}\right)t\right]}} \leq \nu^{(1)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(1)} + (\bar{C})^{(1)}(\bar{\nu}_2)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}}{1 + (\bar{C})^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}} \leq (\bar{\nu}_1)^{(1)}$$

If $0 < (\nu_1)^{(1)} \leq (\bar{\nu}_1)^{(1)} \leq \boxed{(\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}}$ , we obtain

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$$(\nu_1)^{(1)} \leq \nu^{(1)}(t) \leq \frac{(\bar{\nu}_1)^{(1)} + (\bar{C})^{(1)}(\bar{\nu}_2)^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}}{1 + (\bar{C})^{(1)} e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)} - (\bar{\nu}_2)^{(1)}\right)t\right]}} \leq (\nu_0)^{(1)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(1)}(t)$ :-

$$(m_2)^{(1)} \leq \nu^{(1)}(t) \leq (m_1)^{(1)}, \quad \boxed{\nu^{(1)}(t) = \frac{G_{13}(t)}{G_{14}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(1)}(t)$ :-

$$(\mu_2)^{(1)} \leq u^{(1)}(t) \leq (\mu_1)^{(1)}, \quad \boxed{u^{(1)}(t) = \frac{T_{13}(t)}{T_{14}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{13}^{''})^{(1)} = (a_{14}^{''})^{(1)}, then\ (\sigma_1)^{(1)} = (\sigma_2)^{(1)}$ and in this case $(\nu_1)^{(1)} = (\bar{\nu}_1)^{(1)}$ if in addition $(\nu_0)^{(1)} = (\nu_1)^{(1)}$ then $\nu^{(1)}(t) = (\nu_0)^{(1)}$ and as a consequence $G_{13}(t) = (\nu_0)^{(1)} G_{14}(t)$ this also defines $(\nu_0)^{(1)}$ for the special case

Analogously if $(b_{13}^{''})^{(1)} = (b_{14}^{''})^{(1)}, then\ (\tau_1)^{(1)} = (\tau_2)^{(1)}$ and then

$(u_1)^{(1)} = (\bar{u}_1)^{(1)}$ if in addition $(u_0)^{(1)} = (u_1)^{(1)}$ then $T_{13}(t) = (u_0)^{(1)} T_{14}(t)$ This is an important consequence of the relation between $(\nu_1)^{(1)}$ and $(\bar{\nu}_1)^{(1)}$, and definition of $(u_0)^{(1)}$.

**Proof** : From global equations we obtain 387

$$\frac{d\nu^{(2)}}{dt} = (a_{16})^{(2)} - \left((a_{16}^{'})^{(2)} - (a_{17}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17}, t)\right) - (a_{17}^{''})^{(2)}(T_{17}, t)\nu^{(2)} - (a_{17})^{(2)}\nu^{(2)}$$

**Definition of** $\nu^{(2)}$ :- $\boxed{\nu^{(2)} = \frac{G_{16}}{G_{17}}}$ 388

It follows 389

$$-\left((a_{17})^{(2)}\left(\nu^{(2)}\right)^2 + (\sigma_2)^{(2)}\nu^{(2)} - (a_{16})^{(2)}\right) \le \frac{d\nu^{(2)}}{dt} \le -\left((a_{17})^{(2)}\left(\nu^{(2)}\right)^2 + (\sigma_1)^{(2)}\nu^{(2)} - (a_{16})^{(2)}\right)$$

From which one obtains 390

**Definition of** $(\bar{\nu}_1)^{(2)}, (\nu_0)^{(2)}$ :-

For $0 < (\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0} < (\nu_1)^{(2)} < (\bar{\nu}_1)^{(2)}$

$$\nu^{(2)}(t) \ge \frac{(\nu_1)^{(2)} + (C)^{(2)}(\nu_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_0)^{(2)}\right)t\right]}}{1 + (C)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_0)^{(2)}\right)t\right]}} \quad , \quad \boxed{(C)^{(2)} = \frac{(\nu_1)^{(2)} - (\nu_0)^{(2)}}{(\nu_0)^{(2)} - (\nu_2)^{(2)}}}$$

it follows $(\nu_0)^{(2)} \le \nu^{(2)}(t) \le (\nu_1)^{(2)}$

In the same manner , we get 391

$$\nu^{(2)}(t) \le \frac{(\bar{\nu}_1)^{(2)} + (\bar{C})^{(2)}(\bar{\nu}_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}}{1 + (\bar{C})^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(2)} = \frac{(\bar{\nu}_1)^{(2)} - (\nu_0)^{(2)}}{(\nu_0)^{(2)} - (\bar{\nu}_2)^{(2)}}}$$

From which we deduce $(\nu_0)^{(2)} \le \nu^{(2)}(t) \le (\bar{\nu}_1)^{(2)}$ 392

If $0 < (\nu_1)^{(2)} < (\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0} < (\bar{\nu}_1)^{(2)}$ we find like in the previous case, 393

$$(\nu_1)^{(2)} \le \frac{(\nu_1)^{(2)} + (C)^{(2)}(\nu_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_2)^{(2)}\right)t\right]}}{1 + (C)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_2)^{(2)}\right)t\right]}} \le \nu^{(2)}(t) \le$$

| | |
|---|---|
| $\frac{(\bar{\nu}_1)^{(2)}+(\bar{C})^{(2)}(\bar{\nu}_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)}-(\bar{\nu}_2)^{(2)}\right)t\right]}}{1+(\bar{C})^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)}-(\bar{\nu}_2)^{(2)}\right)t\right]}} \leq (\bar{\nu}_1)^{(2)}$ | |
| If $0 < (\nu_1)^{(2)} \leq (\bar{\nu}_1)^{(2)} \leq (\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0}$ , we obtain<br><br>$(\nu_1)^{(2)} \leq \nu^{(2)}(t) \leq \frac{(\bar{\nu}_1)^{(2)}+(\bar{C})^{(2)}(\bar{\nu}_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)}-(\bar{\nu}_2)^{(2)}\right)t\right]}}{1+(\bar{C})^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)}-(\bar{\nu}_2)^{(2)}\right)t\right]}} \leq (\nu_0)^{(2)}$<br><br>And so with the notation of the first part of condition (c) , we have | 394 |
| **Definition of** $\nu^{(2)}(t)$ :-<br><br>$(m_2)^{(2)} \leq \nu^{(2)}(t) \leq (m_1)^{(2)}$, $\boxed{\nu^{(2)}(t) = \frac{G_{16}(t)}{G_{17}(t)}}$ | 395 |
| In a completely analogous way, we obtain<br><br>**Definition of** $u^{(2)}(t)$ :-<br><br>$(\mu_2)^{(2)} \leq u^{(2)}(t) \leq (\mu_1)^{(2)}$, $\boxed{u^{(2)}(t) = \frac{T_{16}(t)}{T_{17}(t)}}$ | 396 |
| Now, using this result and replacing it in global equations we get easily the result stated in the theorem. | |
| **Particular case :**<br><br>If $(a_{16}^{''})^{(2)} = (a_{17}^{''})^{(2)}, then\ (\sigma_1)^{(2)} = (\sigma_2)^{(2)}$ and in this case $(\nu_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ if in addition $(\nu_0)^{(2)} = (\nu_1)^{(2)}$ then $\nu^{(2)}(t) = (\nu_0)^{(2)}$ and as a consequence $G_{16}(t) = (\nu_0)^{(2)}G_{17}(t)$<br><br>Analogously if $(b_{16}^{''})^{(2)} = (b_{17}^{''})^{(2)}, then\ (\tau_1)^{(2)} = (\tau_2)^{(2)}$ and then<br><br>$(u_1)^{(2)} = (\bar{u}_1)^{(2)}$if in addition $(u_0)^{(2)} = (u_1)^{(2)}$ then $T_{16}(t) = (u_0)^{(2)}T_{17}(t)$ This is an important consequence of the relation between $(\nu_1)^{(2)}$ and $(\bar{\nu}_1)^{(2)}$ | 397 |
| **Proof** : From global equations we obtain<br><br>$\frac{d\nu^{(3)}}{dt} = (a_{20})^{(3)} - \left((a_{20}^{'})^{(3)} - (a_{21}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21},t)\right) - (a_{21}^{''})^{(3)}(T_{21},t)\nu^{(3)} - (a_{21})^{(3)}\nu^{(3)}$ | 398 |
| **Definition of** $\nu^{(3)}$ :- $\boxed{\nu^{(3)} = \frac{G_{20}}{G_{21}}}$<br><br>It follows<br><br>$-\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_2)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right) \leq \frac{d\nu^{(3)}}{dt} \leq -\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_1)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right)$ | 399 |
| | 400 |

From which one obtains

For $0 < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\nu_1)^{(3)} < (\bar{\nu}_1)^{(3)}$

$$\nu^{(3)}(t) \geq \frac{(\nu_1)^{(3)}+(C)^{(3)}(\nu_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_0)^{(3)}\right)t\right]}}{1+(C)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_0)^{(3)}\right)t\right]}} \quad , \quad \boxed{(C)^{(3)} = \frac{(\nu_1)^{(3)}-(\nu_0)^{(3)}}{(\nu_0)^{(3)}-(\nu_2)^{(3)}}}$$

it follows $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\nu_1)^{(3)}$

In the same manner , we get

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$$\nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(3)} = \frac{(\bar{\nu}_1)^{(3)}-(\nu_0)^{(3)}}{(\nu_0)^{(3)}-(\bar{\nu}_2)^{(3)}}}$$

**Definition of** $(\bar{\nu}_1)^{(3)}$ :-

From which we deduce $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\bar{\nu}_1)^{(3)}$

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If $0 < (\nu_1)^{(3)} < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\bar{\nu}_1)^{(3)}$ we find like in the previous case,

$$(\nu_1)^{(3)} \leq \frac{(\nu_1)^{(3)}+(C)^{(3)}(\nu_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}}{1+(C)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}} \leq \nu^{(3)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\bar{\nu}_1)^{(3)}$$

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If $0 < (\nu_1)^{(3)} \leq (\bar{\nu}_1)^{(3)} \leq (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0}$ , we obtain

$$(\nu_1)^{(3)} \leq \nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\nu_0)^{(3)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(3)}(t)$ :-

$$(m_2)^{(3)} \leq \nu^{(3)}(t) \leq (m_1)^{(3)}, \quad \boxed{\nu^{(3)}(t) = \frac{G_{20}(t)}{G_{21}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(3)}(t)$ :-

$$(\mu_2)^{(3)} \leq u^{(3)}(t) \leq (\mu_1)^{(3)}, \quad \boxed{u^{(3)}(t) = \frac{T_{20}(t)}{T_{21}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{20}^{''})^{(3)} = (a_{21}^{''})^{(3)}, then\ (\sigma_1)^{(3)} = (\sigma_2)^{(3)}$ and in this case $(\nu_1)^{(3)} = (\bar{\nu}_1)^{(3)}$ if in addition $(\nu_0)^{(3)} = (\nu_1)^{(3)}$ then $\nu^{(3)}(t) = (\nu_0)^{(3)}$ and as a consequence $G_{20}(t) = (\nu_0)^{(3)} G_{21}(t)$

Analogously if $(b_{20}^{''})^{(3)} = (b_{21}^{''})^{(3)}, then\ (\tau_1)^{(3)} = (\tau_2)^{(3)}$ and then

$(u_1)^{(3)} = (\bar{u}_1)^{(3)}$ if in addition $(u_0)^{(3)} = (u_1)^{(3)}$ then $T_{20}(t) = (u_0)^{(3)} T_{21}(t)$ This is an important consequence of the relation between $(\nu_1)^{(3)}$ and $(\bar{\nu}_1)^{(3)}$

**Proof** : From global equations we obtain

$$\frac{d\nu^{(4)}}{dt} = (a_{24})^{(4)} - \left((a_{24}^{'})^{(4)} - (a_{25}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25}, t)\right) - (a_{25}^{''})^{(4)}(T_{25}, t)\nu^{(4)} - (a_{25})^{(4)}\nu^{(4)}$$

**Definition of** $\nu^{(4)}$ :- $\boxed{\nu^{(4)} = \frac{G_{24}}{G_{25}}}$

It follows

$$-\left((a_{25})^{(4)}\left(\nu^{(4)}\right)^2 + (\sigma_2)^{(4)}\nu^{(4)} - (a_{24})^{(4)}\right) \leq \frac{d\nu^{(4)}}{dt} \leq -\left((a_{25})^{(4)}\left(\nu^{(4)}\right)^2 + (\sigma_4)^{(4)}\nu^{(4)} - (a_{24})^{(4)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(4)}, (\nu_0)^{(4)}$ :-

For $0 < \boxed{(\nu_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0}} < (\nu_1)^{(4)} < (\bar{\nu}_1)^{(4)}$

$$\nu^{(4)}(t) \geq \frac{(\nu_1)^{(4)} + (C)^{(4)}(\nu_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_0)^{(4)}\right)t\right]}}{4 + (C)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_0)^{(4)}\right)t\right]}} \quad , \quad \boxed{(C)^{(4)} = \frac{(\nu_1)^{(4)} - (\nu_0)^{(4)}}{(\nu_0)^{(4)} - (\nu_2)^{(4)}}}$$

it follows $(\nu_0)^{(4)} \leq \nu^{(4)}(t) \leq (\nu_1)^{(4)}$

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In the same manner , we get

$$\nu^{(4)}(t) \leq \frac{(\bar{\nu}_1)^{(4)} + (\bar{C})^{(4)}(\bar{\nu}_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}}{4 + (\bar{C})^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(4)} = \frac{(\bar{\nu}_1)^{(4)} - (\nu_0)^{(4)}}{(\nu_0)^{(4)} - (\bar{\nu}_2)^{(4)}}}$$

From which we deduce $(\nu_0)^{(4)} \leq \nu^{(4)}(t) \leq (\bar{\nu}_1)^{(4)}$

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If $0 < (\nu_1)^{(4)} < (\nu_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0} < (\bar{\nu}_1)^{(4)}$ we find like in the previous case,

$$(\nu_1)^{(4)} \leq \frac{(\nu_1)^{(4)} + (C)^{(4)}(\nu_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_2)^{(4)}\right)t\right]}}{1 + (C)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_2)^{(4)}\right)t\right]}} \leq \nu^{(4)}(t) \leq$$

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$$\frac{(\bar{\nu}_1)^{(4)}+(\bar{C})^{(4)}(\bar{\nu}_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)}-(\bar{\nu}_2)^{(4)}\right)t\right]}}{1+(\bar{C})^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)}-(\bar{\nu}_2)^{(4)}\right)t\right]}} \leq (\bar{\nu}_1)^{(4)}$$

If $0 < (\nu_1)^{(4)} \leq (\bar{\nu}_1)^{(4)} \leq \boxed{(\nu_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0}}$ , we obtain

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$$(\nu_1)^{(4)} \leq \nu^{(4)}(t) \leq \frac{(\bar{\nu}_1)^{(4)}+(\bar{C})^{(4)}(\bar{\nu}_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)}-(\bar{\nu}_2)^{(4)}\right)t\right]}}{1+(\bar{C})^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)}-(\bar{\nu}_2)^{(4)}\right)t\right]}} \leq (\nu_0)^{(4)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(4)}(t)$ :-

$$(m_2)^{(4)} \leq \nu^{(4)}(t) \leq (m_1)^{(4)}, \quad \boxed{\nu^{(4)}(t) = \frac{G_{24}(t)}{G_{25}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(4)}(t)$ :-

$$(\mu_2)^{(4)} \leq u^{(4)}(t) \leq (\mu_1)^{(4)}, \quad \boxed{u^{(4)}(t) = \frac{T_{24}(t)}{T_{25}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{24}^{''})^{(4)} = (a_{25}^{''})^{(4)}, then\ (\sigma_1)^{(4)} = (\sigma_2)^{(4)}$ and in this case $(\nu_1)^{(4)} = (\bar{\nu}_1)^{(4)}$ if in addition $(\nu_0)^{(4)} = (\nu_1)^{(4)}$ then $\nu^{(4)}(t) = (\nu_0)^{(4)}$ and as a consequence $G_{24}(t) = (\nu_0)^{(4)}G_{25}(t)$ **this also defines** $(\nu_0)^{(4)}$ **for the special case** .

Analogously if $(b_{24}^{''})^{(4)} = (b_{25}^{''})^{(4)}, then\ (\tau_1)^{(4)} = (\tau_2)^{(4)}$ and then $(u_1)^{(4)} = (\bar{u}_4)^{(4)}$ if in addition $(u_0)^{(4)} = (u_1)^{(4)}$ then $T_{24}(t) = (u_0)^{(4)}T_{25}(t)$ This is an important consequence of the relation between $(\nu_1)^{(4)}$ and $(\bar{\nu}_1)^{(4)}$, **and definition of** $(u_0)^{(4)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(5)}}{dt} = (a_{28})^{(5)} - \left((a_{28}^{'})^{(5)} - (a_{29}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}, t)\right) - (a_{29}^{''})^{(5)}(T_{29}, t)\nu^{(5)} - (a_{29})^{(5)}\nu^{(5)}$$

**Definition of** $\nu^{(5)}$ :- $\boxed{\nu^{(5)} = \frac{G_{28}}{G_{29}}}$

It follows

$$-\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_2)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right) \leq \frac{d\nu^{(5)}}{dt} \leq -\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_1)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(5)}, (\nu_0)^{(5)}$ :-

For $0 < \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}} < (\nu_1)^{(5)} < (\bar{\nu}_1)^{(5)}$

$$\nu^{(5)}(t) \geq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}}{5 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}} \quad , \quad \boxed{(C)^{(5)} = \frac{(\nu_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\nu_2)^{(5)}}}$$

it follows $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\nu_1)^{(5)}$

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In the same manner , we get

$$\nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)} + (\bar{C})^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{5 + (\bar{C})^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(5)} = \frac{(\bar{\nu}_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\bar{\nu}_2)^{(5)}}}$$

From which we deduce $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\bar{\nu}_5)^{(5)}$

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If $0 < (\nu_1)^{(5)} < (\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0} < (\bar{\nu}_1)^{(5)}$ we find like in the previous case,

$$(\nu_1)^{(5)} \leq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}}{1 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}} \leq \nu^{(5)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(5)} + (C)^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{1 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\bar{\nu}_1)^{(5)}$$

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If $0 < (\nu_1)^{(5)} \leq (\bar{\nu}_1)^{(5)} \leq \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}}$ , we obtain

$$(\nu_1)^{(5)} \leq \nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)} + (C)^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{1 + (\bar{C})^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\nu_0)^{(5)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(5)}(t)$ :-

$$(m_2)^{(5)} \leq \nu^{(5)}(t) \leq (m_1)^{(5)}, \quad \boxed{\nu^{(5)}(t) = \frac{G_{28}(t)}{G_{29}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(5)}(t)$ :-

$$(\mu_2)^{(5)} \leq u^{(5)}(t) \leq (\mu_1)^{(5)}, \quad \boxed{u^{(5)}(t) = \frac{T_{28}(t)}{T_{29}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{28}'')^{(5)} = (a_{29}'')^{(5)}, then\ (\sigma_1)^{(5)} = (\sigma_2)^{(5)}$ and in this case $(\nu_1)^{(5)} = (\bar{\nu}_1)^{(5)}$ if in addition $(\nu_0)^{(5)} = (\nu_5)^{(5)}$ then $\nu^{(5)}(t) = (\nu_0)^{(5)}$ and as a consequence $G_{28}(t) = (\nu_0)^{(5)} G_{29}(t)$ **this also defines** $(\nu_0)^{(5)}$ **for the special case** .

Analogously if $(b_{28}'')^{(5)} = (b_{29}'')^{(5)}, then\ (\tau_1)^{(5)} = (\tau_2)^{(5)}$ and then

$(u_1)^{(5)} = (\bar{u}_1)^{(5)}$if in addition $(u_0)^{(5)} = (u_1)^{(5)}$ then $T_{28}(t) = (u_0)^{(5)}T_{29}(t)$ This is an important consequence of the relation between $(\nu_1)^{(5)}$ and $(\bar{\nu}_1)^{(5)}$, **and definition of** $(u_0)^{(5)}$.

**Proof** : From global equations we obtain

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$$\frac{d\nu^{(6)}}{dt} = (a_{32})^{(6)} - \left((a'_{32})^{(6)} - (a'_{33})^{(6)} + (a''_{32})^{(6)}(T_{33},t)\right) - (a''_{33})^{(6)}(T_{33},t)\nu^{(6)} - (a_{33})^{(6)}\nu^{(6)}$$

**Definition of** $\nu^{(6)}$ :- $\boxed{\nu^{(6)} = \frac{G_{32}}{G_{33}}}$

It follows

$$-\left((a_{33})^{(6)}\left(\nu^{(6)}\right)^2 + (\sigma_2)^{(6)}\nu^{(6)} - (a_{32})^{(6)}\right) \le \frac{d\nu^{(6)}}{dt} \le -\left((a_{33})^{(6)}\left(\nu^{(6)}\right)^2 + (\sigma_1)^{(6)}\nu^{(6)} - (a_{32})^{(6)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(6)}, (\nu_0)^{(6)}$ :-

For $0 < \boxed{(\nu_0)^{(6)} = \frac{G^0_{32}}{G^0_{33}}} < (\nu_1)^{(6)} < (\bar{\nu}_1)^{(6)}$

$$\nu^{(6)}(t) \ge \frac{(\nu_1)^{(6)} + (C)^{(6)}(\nu_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)} - (\nu_0)^{(6)}\right)t\right]}}{1 + (C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)} - (\nu_0)^{(6)}\right)t\right]}} \quad , \quad \boxed{(C)^{(6)} = \frac{(\nu_1)^{(6)} - (\nu_0)^{(6)}}{(\nu_0)^{(6)} - (\nu_2)^{(6)}}}$$

it follows $(\nu_0)^{(6)} \le \nu^{(6)}(t) \le (\nu_1)^{(6)}$

In the same manner , we get

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$$\nu^{(6)}(t) \le \frac{(\bar{\nu}_1)^{(6)} + (\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)} - (\bar{\nu}_2)^{(6)}\right)t\right]}}{1 + (\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)} - (\bar{\nu}_2)^{(6)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(6)} = \frac{(\bar{\nu}_1)^{(6)} - (\nu_0)^{(6)}}{(\nu_0)^{(6)} - (\bar{\nu}_2)^{(6)}}}$$

From which we deduce $(\nu_0)^{(6)} \le \nu^{(6)}(t) \le (\bar{\nu}_1)^{(6)}$

If $0 < (\nu_1)^{(6)} < (\nu_0)^{(6)} = \frac{G^0_{32}}{G^0_{33}} < (\bar{\nu}_1)^{(6)}$ we find like in the previous case,

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$$(\nu_1)^{(6)} \le \frac{(\nu_1)^{(6)} + (C)^{(6)}(\nu_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)} - (\nu_2)^{(6)}\right)t\right]}}{1 + (C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)} - (\nu_2)^{(6)}\right)t\right]}} \le \nu^{(6)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(6)} + (\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)} - (\bar{\nu}_2)^{(6)}\right)t\right]}}{1 + (\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)} - (\bar{\nu}_2)^{(6)}\right)t\right]}} \le (\bar{\nu}_1)^{(6)}$$

If $0 < (\nu_1)^{(6)} \le (\bar{\nu}_1)^{(6)} \le \boxed{(\nu_0)^{(6)} = \frac{G^0_{32}}{G^0_{33}}}$ , we obtain

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$$(\nu_1)^{(6)} \le \nu^{(6)}(t) \le \frac{(\bar{\nu}_1)^{(6)} + (\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)} - (\bar{\nu}_2)^{(6)}\right)t\right]}}{1 + (C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)} - (\bar{\nu}_2)^{(6)}\right)t\right]}} \le (\nu_0)^{(6)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(6)}(t)$ :-

$(m_2)^{(6)} \leq \nu^{(6)}(t) \leq (m_1)^{(6)}$, $\boxed{\nu^{(6)}(t) = \frac{G_{32}(t)}{G_{33}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(6)}(t)$ :-

$(\mu_2)^{(6)} \leq u^{(6)}(t) \leq (\mu_1)^{(6)}$, $\boxed{u^{(6)}(t) = \frac{T_{32}(t)}{T_{33}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{32}^{''})^{(6)} = (a_{33}^{''})^{(6)}, then\ (\sigma_1)^{(6)} = (\sigma_2)^{(6)}$ and in this case $(\nu_1)^{(6)} = (\bar{\nu}_1)^{(6)}$ if in addition $(\nu_0)^{(6)} = (\nu_1)^{(6)}$ then $\nu^{(6)}(t) = (\nu_0)^{(6)}$ and as a consequence $G_{32}(t) = (\nu_0)^{(6)} G_{33}(t)$ **this also defines** $(\nu_0)^{(6)}$ **for the special case** .

Analogously if $(b_{32}^{''})^{(6)} = (b_{33}^{''})^{(6)}, then\ (\tau_1)^{(6)} = (\tau_2)^{(6)}$ and then $(u_1)^{(6)} = (\bar{u}_1)^{(6)}$ if in addition $(u_0)^{(6)} = (u_1)^{(6)}$ then $T_{32}(t) = (u_0)^{(6)} T_{33}(t)$ This is an important consequence of the relation between $(\nu_1)^{(6)}$ and $(\bar{\nu}_1)^{(6)}$, **and definition of** $(u_0)^{(6)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(7)}}{dt} = (a_{36})^{(7)} - \left((a_{36}^{'})^{(7)} - (a_{37}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37}, t)\right) - (a_{37}^{''})^{(7)}(T_{37}, t)\nu^{(7)} - (a_{37})^{(7)}\nu^{(7)}$$

**Definition of** $\nu^{(7)}$ :- $\boxed{\nu^{(7)} = \frac{G_{36}}{G_{37}}}$

It follows

$$-\left((a_{37})^{(7)}\left(\nu^{(7)}\right)^2 + (\sigma_2)^{(7)}\nu^{(7)} - (a_{36})^{(7)}\right) \leq \frac{d\nu^{(7)}}{dt} \leq -\left((a_{37})^{(7)}\left(\nu^{(7)}\right)^2 + (\sigma_1)^{(7)}\nu^{(7)} - (a_{36})^{(7)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(7)}, (\nu_0)^{(7)}$ :-

For $0 < \boxed{(\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0}} < (\nu_1)^{(7)} < (\bar{\nu}_1)^{(7)}$

$$\nu^{(7)}(t) \geq \frac{(\nu_1)^{(7)} + (C)^{(7)}(\nu_2)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)} - (\nu_0)^{(7)}\right)t\right]}}{1 + (C)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)} - (\nu_0)^{(7)}\right)t\right]}} \quad , \quad \boxed{(C)^{(7)} = \frac{(\nu_1)^{(7)} - (\nu_0)^{(7)}}{(\nu_0)^{(7)} - (\nu_2)^{(7)}}}$$

it follows $(\nu_0)^{(7)} \leq \nu^{(7)}(t) \leq (\nu_1)^{(7)}$

In the same manner , we get

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$$\nu^{(7)}(t) \le \frac{(\bar{\nu}_1)^{(7)}+(\bar{C})^{(7)}(\bar{\nu}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}}{1+(\bar{C})^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(7)} = \frac{(\bar{\nu}_1)^{(7)}-(\nu_0)^{(7)}}{(\nu_0)^{(7)}-(\bar{\nu}_2)^{(7)}}}$$

From which we deduce $(\nu_0)^{(7)} \le \nu^{(7)}(t) \le (\bar{\nu}_1)^{(7)}$

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If $0 < (\nu_1)^{(7)} < (\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0} < (\bar{\nu}_1)^{(7)}$ we find like in the previous case,

$$(\nu_1)^{(7)} \le \frac{(\nu_1)^{(7)}+(C)^{(7)}(\nu_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)}-(\nu_2)^{(7)}\right)t\right]}}{1+(C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)}-(\nu_2)^{(7)}\right)t\right]}} \le \nu^{(7)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(7)}+(\bar{C})^{(7)}(\bar{\nu}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}}{1+(\bar{C})^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}} \le (\bar{\nu}_1)^{(7)}$$

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If $0 < (\nu_1)^{(7)} \le (\bar{\nu}_1)^{(7)} \le \boxed{(\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0}}$ , we obtain

$$(\nu_1)^{(7)} \le \nu^{(7)}(t) \le \frac{(\bar{\nu}_1)^{(7)}+(C)^{(7)}(\bar{\nu}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}}{1+(C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)}-(\bar{\nu}_2)^{(7)}\right)t\right]}} \le (\nu_0)^{(7)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(7)}(t)$ :-

$$(m_2)^{(7)} \le \nu^{(7)}(t) \le (m_1)^{(7)}, \quad \boxed{\nu^{(7)}(t) = \frac{G_{36}(t)}{G_{37}(t)}}$$

In a completely analogous way, we obtain

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**Definition of** $u^{(7)}(t)$ :-

$$(\mu_2)^{(7)} \le u^{(7)}(t) \le (\mu_1)^{(7)}, \quad \boxed{u^{(7)}(t) = \frac{T_{36}(t)}{T_{37}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{36}'')^{(7)} = (a_{37}'')^{(7)}, then\ (\sigma_1)^{(7)} = (\sigma_2)^{(7)}$ and in this case $(\nu_1)^{(7)} = (\bar{\nu}_1)^{(7)}$ if in addition $(\nu_0)^{(7)} = (\nu_1)^{(7)}$ then $\nu^{(7)}(t) = (\nu_0)^{(7)}$ and as a consequence $G_{36}(t) = (\nu_0)^{(7)}G_{37}(t)$ **this also defines** $(\nu_0)^{(7)}$ **for the special case** .

Analogously if $(b_{36}'')^{(7)} = (b_{37}'')^{(7)}, then\ (\tau_1)^{(7)} = (\tau_2)^{(7)}$ and then $(u_1)^{(7)} = (\bar{u}_1)^{(7)}$if in addition $(u_0)^{(7)} = (u_1)^{(7)}$ then $T_{36}(t) = (u_0)^{(7)}T_{37}(t)$ This is an important consequence of the relation between $(\nu_1)^{(7)}$ and $(\bar{\nu}_1)^{(7)}$, **and definition of** $(u_0)^{(7)}$.

**Proof** : From global equations we obtain 421

$$\frac{dv^{(8)}}{dt} = (a_{40})^{(8)} - \left((a'_{40})^{(8)} - (a'_{41})^{(8)} + (a''_{40})^{(8)}(T_{41},t)\right) - (a''_{41})^{(8)}(T_{41},t)v^{(8)} - (a_{41})^{(8)}v^{(8)}$$

**Definition of** $v^{(8)}$ :- $\boxed{v^{(8)} = \frac{G_{40}}{G_{41}}}$

It follows

$$-\left((a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_2)^{(8)}v^{(8)} - (a_{40})^{(8)}\right) \le \frac{dv^{(8)}}{dt} \le -\left((a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_1)^{(8)}v^{(8)} - (a_{40})^{(8)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(8)}, (\nu_0)^{(8)}$ :-

For $0 < \boxed{(\nu_0)^{(8)} = \frac{G^0_{40}}{G^0_{41}}} < (\nu_1)^{(8)} < (\bar{\nu}_1)^{(8)}$

$$\nu^{(8)}(t) \ge \frac{(\nu_1)^{(8)} + (C)^{(8)}(\nu_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}}{1 + (C)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}} \quad , \quad \boxed{(C)^{(8)} = \frac{(\nu_1)^{(8)} - (\nu_0)^{(8)}}{(\nu_0)^{(8)} - (\nu_2)^{(8)}}}$$

it follows $(\nu_0)^{(8)} \le \nu^{(8)}(t) \le (\nu_1)^{(8)}$

In the same manner , we get 422

$$\nu^{(8)}(t) \le \frac{(\bar{\nu}_1)^{(8)} + (\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}}{1 + (\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(8)} = \frac{(\bar{\nu}_1)^{(8)} - (\nu_0)^{(8)}}{(\nu_0)^{(8)} - (\bar{\nu}_2)^{(8)}}}$$

From which we deduce $(\nu_0)^{(8)} \le \nu^{(8)}(t) \le (\bar{\nu}_8)^{(8)}$

If $0 < (\nu_1)^{(8)} < (\nu_0)^{(8)} = \frac{G^0_{40}}{G^0_{41}} < (\bar{\nu}_1)^{(8)}$ we find like in the previous case, 423

$$(\nu_1)^{(8)} \le \frac{(\nu_1)^{(8)} + (C)^{(8)}(\nu_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_2)^{(8)}\right)t\right]}}{1 + (C)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_2)^{(8)}\right)t\right]}} \le \nu^{(8)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(8)} + (C)^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}}{1 + (\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}} \le (\bar{\nu}_1)^{(8)}$$

If $0 < (\nu_1)^{(8)} \le (\bar{\nu}_1)^{(8)} \le \boxed{(\nu_0)^{(8)} = \frac{G^0_{40}}{G^0_{41}}}$ , we obtain 424

$$(\nu_1)^{(8)} \le \nu^{(8)}(t) \le \frac{(\bar{\nu}_1)^{(8)} + (\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}}{1 + (\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)} - (\bar{\nu}_2)^{(8)}\right)t\right]}} \le (\nu_0)^{(8)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $v^{(8)}(t)$ :-

$(m_2)^{(8)} \leq v^{(8)}(t) \leq (m_1)^{(8)}$, $\boxed{v^{(8)}(t) = \frac{G_{40}(t)}{G_{41}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(8)}(t)$ :-

$(\mu_2)^{(8)} \leq u^{(8)}(t) \leq (\mu_1)^{(8)}$, $\boxed{u^{(8)}(t) = \frac{T_{40}(t)}{T_{41}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{40}^{''})^{(8)} = (a_{41}^{''})^{(8)}, then\ (\sigma_1)^{(8)} = (\sigma_2)^{(8)}$ and in this case $(\nu_1)^{(8)} = (\bar{\nu}_1)^{(8)}$ if in addition $(\nu_0)^{(8)} = (\nu_1)^{(8)}$ then $v^{(8)}(t) = (\nu_0)^{(8)}$ and as a consequence $G_{40}(t) = (\nu_0)^{(8)} G_{41}(t)$ **this also defines** $(\nu_0)^{(8)}$ **for the special case** .

Analogously if $(b_{40}^{''})^{(8)} = (b_{41}^{''})^{(8)}, then\ (\tau_1)^{(8)} = (\tau_2)^{(8)}$ and then
$(u_1)^{(8)} = (\bar{u}_1)^{(8)}$if in addition $(u_0)^{(8)} = (u_1)^{(8)}$ then $T_{40}(t) = (u_0)^{(8)} T_{41}(t)$ This is an important consequence of the relation between $(\nu_1)^{(8)}$ and $(\bar{\nu}_1)^{(8)}$, **and definition of** $(u_0)^{(8)}$.

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**Proof** : From 99,20,44,22,23,44 we obtain

$$\frac{dv^{(9)}}{dt} = (a_{44})^{(9)} - \left((a_{44}^{'})^{(9)} - (a_{45}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45}, t)\right) - (a_{45}^{''})^{(9)}(T_{45}, t)v^{(9)} - (a_{45})^{(9)}v^{(9)}$$

**Definition of** $v^{(9)}$ :- $\boxed{v^{(9)} = \frac{G_{44}}{G_{45}}}$

It follows

$$-\left((a_{45})^{(9)}\left(v^{(9)}\right)^2 + (\sigma_2)^{(9)}v^{(9)} - (a_{44})^{(9)}\right) \leq \frac{dv^{(9)}}{dt} \leq -\left((a_{45})^{(9)}\left(v^{(9)}\right)^2 + (\sigma_1)^{(9)}v^{(9)} - (a_{44})^{(9)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(9)}, (\nu_0)^{(9)}$ :-

For $0 < \boxed{(\nu_0)^{(9)} = \frac{G_{44}^0}{G_{45}^0}} < (\nu_1)^{(9)} < (\bar{\nu}_1)^{(9)}$

$$v^{(9)}(t) \geq \frac{(\nu_1)^{(9)} + (C)^{(9)}(\nu_2)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_0)^{(9)}\right)t\right]}}{1 + (C)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_0)^{(9)}\right)t\right]}} \quad , \quad \boxed{(C)^{(9)} = \frac{(\nu_1)^{(9)} - (\nu_0)^{(9)}}{(\nu_0)^{(9)} - (\nu_2)^{(9)}}}$$

it follows $(\nu_0)^{(9)} \leq v^{(9)}(t) \leq (\nu_9)^{(9)}$

In the same manner , we get

$$\nu^{(9)}(t) \le \frac{(\bar{\nu}_1)^{(9)}+(\bar{C})^{(9)}(\bar{\nu}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)}-(\bar{\nu}_2)^{(9)}\right)t\right]}}{1+(\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)}-(\bar{\nu}_2)^{(9)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(9)} = \frac{(\bar{\nu}_1)^{(9)}-(\nu_0)^{(9)}}{(\nu_0)^{(9)}-(\bar{\nu}_2)^{(9)}}}$$

From which we deduce $(\nu_0)^{(9)} \le \nu^{(9)}(t) \le (\bar{\nu}_1)^{(9)}$

If $0 < (\nu_1)^{(9)} < (\nu_0)^{(9)} = \frac{G_{44}^0}{G_{45}^0} < (\bar{\nu}_1)^{(9)}$ we find like in the previous case,

$$(\nu_1)^{(9)} \le \frac{(\nu_1)^{(9)}+(C)^{(9)}(\nu_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)}-(\nu_2)^{(9)}\right)t\right]}}{1+(C)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)}-(\nu_2)^{(9)}\right)t\right]}} \le \nu^{(9)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(9)}+(\bar{C})^{(9)}(\bar{\nu}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)}-(\bar{\nu}_2)^{(9)}\right)t\right]}}{1+(\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)}-(\bar{\nu}_2)^{(9)}\right)t\right]}} \le (\bar{\nu}_1)^{(9)}$$

If $0 < (\nu_1)^{(9)} \le (\bar{\nu}_1)^{(9)} \le \boxed{(\nu_0)^{(9)} = \frac{G_{44}^0}{G_{45}^0}}$, we obtain

$$(\nu_1)^{(9)} \le \nu^{(9)}(t) \le \frac{(\bar{\nu}_1)^{(9)}+(\bar{C})^{(9)}(\bar{\nu}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)}-(\bar{\nu}_2)^{(9)}\right)t\right]}}{1+(\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)}-(\bar{\nu}_2)^{(9)}\right)t\right]}} \le (\nu_0)^{(9)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(9)}(t)$ :-

$$(m_2)^{(9)} \le \nu^{(9)}(t) \le (m_1)^{(9)}, \quad \boxed{\nu^{(9)}(t) = \frac{G_{44}(t)}{G_{45}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(9)}(t)$ :-

$$(\mu_2)^{(9)} \le u^{(9)}(t) \le (\mu_1)^{(9)}, \quad \boxed{u^{(9)}(t) = \frac{T_{44}(t)}{T_{45}(t)}}$$

Now, using this result and replacing it in 99, 20,44,22,23, and 44 we get easily the result stated in the theorem.

**Particular case :**

If $(a_{44}^{''})^{(9)} = (a_{45}^{''})^{(9)}, then\ (\sigma_1)^{(9)} = (\sigma_2)^{(9)}$ and in this case $(\nu_1)^{(9)} = (\bar{\nu}_1)^{(9)}$ if in addition $(\nu_0)^{(9)} = (\nu_1)^{(9)}$ then $\nu^{(9)}(t) = (\nu_0)^{(9)}$ and as a consequence $G_{44}(t) = (\nu_0)^{(9)}G_{45}(t)$ **this also defines** $(\nu_0)^{(9)}$ **for the special case** .

Analogously if $(b_{44}^{''})^{(9)} = (b_{45}^{''})^{(9)}, then\ (\tau_1)^{(9)} = (\tau_2)^{(9)}$ and then $(u_1)^{(9)} = (\bar{u}_1)^{(9)}$ if in addition $(u_0)^{(9)} = (u_1)^{(9)}$ then $T_{44}(t) = (u_0)^{(9)}T_{45}(t)$ This is an important consequence of the relation between $(\nu_1)^{(9)}$ and $(\bar{\nu}_1)^{(9)}$, **and definition of** $(u_0)^{(9)}$.

We can prove the following

**Theorem :** If $(a_i^{''})^{(1)} and\ (b_i^{''})^{(1)}$ are independent on $t$ , and the conditions with the notations

$$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} < 0$$

$$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a_{13})^{(1)}(p_{13})^{(1)} + (a_{14}^{'})^{(1)}(p_{14})^{(1)} + (p_{13})^{(1)}(p_{14})^{(1)} > 0$$

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| $(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} > 0$ ,<br>$(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} - (b_{13}^{'})^{(1)}(r_{14})^{(1)} - (b_{14}^{'})^{(1)}(r_{14})^{(1)} + (r_{13})^{(1)}(r_{14})^{(1)} < 0$<br>$with\ \ (p_{13})^{(1)}, (r_{14})^{(1)}$ as defined by equation are satisfied , then the system | |
| **Theorem :** If $(a_i^{''})^{(2)} and\ (b_i^{''})^{(2)}$ are independent on t , and the conditions with the notations | 426 |
| $(a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} < 0$ | 427 |
| $(a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a_{16})^{(2)}(p_{16})^{(2)} + (a_{17}^{'})^{(2)}(p_{17})^{(2)} + (p_{16})^{(2)}(p_{17})^{(2)} > 0$ | 428 |
| $(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} > 0$ , | 429 |
| $(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} - (b_{16}^{'})^{(2)}(r_{17})^{(2)} - (b_{17}^{'})^{(2)}(r_{17})^{(2)} + (r_{16})^{(2)}(r_{17})^{(2)} < 0$<br>$with\ \ (p_{16})^{(2)}, (r_{17})^{(2)}$ as defined by equation are satisfied , then the system | 430 |
| **Theorem :** If $(a_i^{''})^{(3)} and\ (b_i^{''})^{(3)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} < 0$<br>$(a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a_{20})^{(3)}(p_{20})^{(3)} + (a_{21}^{'})^{(3)}(p_{21})^{(3)} + (p_{20})^{(3)}(p_{21})^{(3)} > 0$<br>$(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} > 0$ ,<br>$(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} - (b_{20}^{'})^{(3)}(r_{21})^{(3)} - (b_{21}^{'})^{(3)}(r_{21})^{(3)} + (r_{20})^{(3)}(r_{21})^{(3)} < 0$<br>$with\ \ (p_{20})^{(3)}, (r_{21})^{(3)}$ as defined by equation are satisfied , then the system | 431 |
| We can prove the following<br>**Theorem :** If $(a_i^{''})^{(4)} and\ (b_i^{''})^{(4)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} < 0$<br>$(a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a_{24})^{(4)}(p_{24})^{(4)} + (a_{25}^{'})^{(4)}(p_{25})^{(4)} + (p_{24})^{(4)}(p_{25})^{(4)} > 0$<br>$(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} > 0$ ,<br>$(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} - (b_{24}^{'})^{(4)}(r_{25})^{(4)} - (b_{25}^{'})^{(4)}(r_{25})^{(4)} + (r_{24})^{(4)}(r_{25})^{(4)} < 0$<br>$with\ \ (p_{24})^{(4)}, (r_{25})^{(4)}$ as defined by equation are satisfied , then the system | 432 |
| **Theorem :** If $(a_i^{''})^{(5)} and\ (b_i^{''})^{(5)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} < 0$<br>$(a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a_{28})^{(5)}(p_{28})^{(5)} + (a_{29}^{'})^{(5)}(p_{29})^{(5)} + (p_{28})^{(5)}(p_{29})^{(5)} > 0$ | 433 |

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| $(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} > 0$ ,<br>$(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} - (b_{28}^{'})^{(5)}(r_{29})^{(5)} - (b_{29}^{'})^{(5)}(r_{29})^{(5)} + (r_{28})^{(5)}(r_{29})^{(5)} < 0$<br>$with\ \ (p_{28})^{(5)}, (r_{29})^{(5)}$ as defined by equation are satisfied , then the system | |
| **Theorem** If $(a_i^{''})^{(6)} and\ (b_i^{''})^{(6)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} < 0$<br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32})^{(6)}(p_{32})^{(6)} + (a_{33}^{'})^{(6)}(p_{33})^{(6)} + (p_{32})^{(6)}(p_{33})^{(6)} > 0$<br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} > 0$ ,<br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} - (b_{32}^{'})^{(6)}(r_{33})^{(6)} - (b_{33}^{'})^{(6)}(r_{33})^{(6)} + (r_{32})^{(6)}(r_{33})^{(6)} < 0$<br>$with\ \ (p_{32})^{(6)}, (r_{33})^{(6)}$ as defined by equation are satisfied , then the system | 434 |
| **Theorem :** If $(a_i^{''})^{(7)} and\ (b_i^{''})^{(7)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} < 0$<br>$(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36})^{(7)}(p_{36})^{(7)} + (a_{37}^{'})^{(7)}(p_{37})^{(7)} + (p_{36})^{(7)}(p_{37})^{(7)} > 0$<br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} > 0$ ,<br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} - (b_{36}^{'})^{(7)}(r_{37})^{(7)} - (b_{37}^{'})^{(7)}(r_{37})^{(7)} + (r_{36})^{(7)}(r_{37})^{(7)} < 0$<br>$with\ \ (p_{36})^{(7)}, (r_{37})^{(7)}$ as defined by equation are satisfied , then the system | 435 |
| **Theorem :** If $(a_i^{''})^{(8)} and\ (b_i^{''})^{(8)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} < 0$<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40})^{(8)}(p_{40})^{(8)} + (a_{41}^{'})^{(8)}(p_{41})^{(8)} + (p_{40})^{(8)}(p_{41})^{(8)} > 0$<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} > 0$ ,<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} - (b_{40}^{'})^{(8)}(r_{41})^{(8)} - (b_{41}^{'})^{(8)}(r_{41})^{(8)} + (r_{40})^{(8)}(r_{41})^{(8)} < 0$<br>$with\ \ (p_{40})^{(8)}, (r_{41})^{(8)}$ as defined by equation are satisfied , then the system | 436 |
| **Theorem :** If $(a_i^{''})^{(9)} and\ (b_i^{''})^{(9)}$ are independent on $t$ , and the conditions (with the notations 45,46,27,28)<br>$(a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} < 0$ | 436 A |

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| $(a'_{44})^{(9)}(a'_{45})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44})^{(9)}(p_{44})^{(9)} + (a'_{45})^{(9)}(p_{45})^{(9)} + (p_{44})^{(9)}(p_{45})^{(9)} > 0$<br>$(b'_{44})^{(9)}(b'_{45})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} > 0$ ,<br>$(b'_{44})^{(9)}(b'_{45})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} - (b'_{44})^{(9)}(r_{45})^{(9)} - (b'_{45})^{(9)}(r_{45})^{(9)} + (r_{44})^{(9)}(r_{45})^{(9)} < 0$<br>$with\ \ (p_{44})^{(9)}, (r_{45})^{(9)}$ as defined by equation 45 are satisfied , then the system | |
| $(a_{13})^{(1)}G_{14} - \left[(a'_{13})^{(1)} + (a''_{13})^{(1)}(T_{14})\right]G_{13} = 0$ | 437 |
| $(a_{14})^{(1)}G_{13} - \left[(a'_{14})^{(1)} + (a''_{14})^{(1)}(T_{14})\right]G_{14} = 0$ | 438 |
| $(a_{15})^{(1)}G_{14} - \left[(a'_{15})^{(1)} + (a''_{15})^{(1)}(T_{14})\right]G_{15} = 0$ | 439 |
| $(b_{13})^{(1)}T_{14} - [(b'_{13})^{(1)} - (b''_{13})^{(1)}(G)\ ]T_{13} = 0$ | 440 |
| $(b_{14})^{(1)}T_{13} - [(b'_{14})^{(1)} - (b''_{14})^{(1)}(G)\ ]T_{14} = 0$ | 441 |
| $(b_{15})^{(1)}T_{14} - [(b'_{15})^{(1)} - (b''_{15})^{(1)}(G)\ ]T_{15} = 0$ | 442 |
| has a unique positive solution , which is an equilibrium solution for the system | |
| $(a_{16})^{(2)}G_{17} - \left[(a'_{16})^{(2)} + (a''_{16})^{(2)}(T_{17})\right]G_{16} = 0$ | 443 |
| $(a_{17})^{(2)}G_{16} - \left[(a'_{17})^{(2)} + (a''_{17})^{(2)}(T_{17})\right]G_{17} = 0$ | 444 |
| $(a_{18})^{(2)}G_{17} - \left[(a'_{18})^{(2)} + (a''_{18})^{(2)}(T_{17})\right]G_{18} = 0$ | 445 |
| $(b_{16})^{(2)}T_{17} - [(b'_{16})^{(2)} - (b''_{16})^{(2)}(G_{19})\ ]T_{16} = 0$ | 446 |
| $(b_{17})^{(2)}T_{16} - [(b'_{17})^{(2)} - (b''_{17})^{(2)}(G_{19})\ ]T_{17} = 0$ | 447 |
| $(b_{18})^{(2)}T_{17} - [(b'_{18})^{(2)} - (b''_{18})^{(2)}(G_{19})\ ]T_{18} = 0$ | 448 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{20})^{(3)}G_{21} - \left[(a'_{20})^{(3)} + (a''_{20})^{(3)}(T_{21})\right]G_{20} = 0$ | 449 |
| $(a_{21})^{(3)}G_{20} - \left[(a'_{21})^{(3)} + (a''_{21})^{(3)}(T_{21})\right]G_{21} = 0$ | 450 |
| $(a_{22})^{(3)}G_{21} - \left[(a'_{22})^{(3)} + (a''_{22})^{(3)}(T_{21})\right]G_{22} = 0$ | 451 |
| $(b_{20})^{(3)}T_{21} - [(b'_{20})^{(3)} - (b''_{20})^{(3)}(G_{23})\ ]T_{20} = 0$ | 452 |
| $(b_{21})^{(3)}T_{20} - [(b'_{21})^{(3)} - (b''_{21})^{(3)}(G_{23})\ ]T_{21} = 0$ | 453 |
| $(b_{22})^{(3)}T_{21} - [(b'_{22})^{(3)} - (b''_{22})^{(3)}(G_{23})\ ]T_{22} = 0$ | 454 |
| has a unique positive solution , which is an equilibrium solution | |

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| $(a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25})\right]G_{24} = 0$ | 455 |
| $(a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25})\right]G_{25} = 0$ | 456 |
| $(a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25})\right]G_{26} = 0$ | 457 |
| $(b_{24})^{(4)}T_{25} - [(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}\left((G_{27})\right)]T_{24} = 0$ | 458 |
| $(b_{25})^{(4)}T_{24} - [(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}\left((G_{27})\right)]T_{25} = 0$ | 459 |
| $(b_{26})^{(4)}T_{25} - [(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}\left((G_{27})\right)]T_{26} = 0$ | 460 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29})\right]G_{28} = 0$ | 461 |
| $(a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29})\right]G_{29} = 0$ | 462 |
| $(a_{30})^{(5)}G_{29} - \left[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29})\right]G_{30} = 0$ | 463 |
| $(b_{28})^{(5)}T_{29} - [(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}(G_{31})]T_{28} = 0$ | 464 |
| $(b_{29})^{(5)}T_{28} - [(b_{29}^{'})^{(5)} - (b_{29}^{''})^{(5)}(G_{31})]T_{29} = 0$ | 465 |
| $(b_{30})^{(5)}T_{29} - [(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}(G_{31})]T_{30} = 0$ | 466 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{32})^{(6)}G_{33} - \left[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33})\right]G_{32} = 0$ | 467 |
| $(a_{33})^{(6)}G_{32} - \left[(a_{33}^{'})^{(6)} + (a_{33}^{''})^{(6)}(T_{33})\right]G_{33} = 0$ | 468 |
| $(a_{34})^{(6)}G_{33} - \left[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33})\right]G_{34} = 0$ | 469 |
| $(b_{32})^{(6)}T_{33} - [(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}(G_{35})]T_{32} = 0$ | 470 |
| $(b_{33})^{(6)}T_{32} - [(b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}(G_{35})]T_{33} = 0$ | 471 |
| $(b_{34})^{(6)}T_{33} - [(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}(G_{35})]T_{34} = 0$ | 472 |
| has a unique positive solution , which is an equilibrium solution | |

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| $(a_{36})^{(7)}G_{37} - \left[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37})\right]G_{36} = 0$ | 473 |
| $(a_{37})^{(7)}G_{36} - \left[(a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37})\right]G_{37} = 0$ | 474 |
| $(a_{38})^{(7)}G_{37} - \left[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37})\right]G_{38} = 0$ | 475 |
| $(b_{36})^{(7)}T_{37} - [(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}(G_{39})]T_{36} = 0$ | 476 |
| $(b_{37})^{(7)}T_{36} - [(b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}(G_{39})\,]T_{37} = 0$ | 477 |
| $(b_{38})^{(7)}T_{37} - [(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}(G_{39})\,]T_{38} = 0$ | 478 |
| $(a_{40})^{(8)}G_{41} - \left[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41})\right]G_{40} = 0$ | 479 |
| $(a_{41})^{(8)}G_{40} - \left[(a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41})\right]G_{41} = 0$ | 480 |
| $(a_{42})^{(8)}G_{41} - \left[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41})\right]G_{42} = 0$ | 481 |
| $(b_{40})^{(8)}T_{41} - [(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}(G_{43})\,]T_{40} = 0$ | 482 |
| $(b_{41})^{(8)}T_{40} - [(b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}(G_{43})\,]T_{41} = 0$ | 483 |
| $(b_{42})^{(8)}T_{41} - [(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}(G_{43})\,]T_{42} = 0$ | 484 |
| $(a_{44})^{(9)}G_{45} - \left[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45})\right]G_{44} = 0$ | 484 A |
| $(a_{45})^{(9)}G_{44} - \left[(a_{45}^{'})^{(9)} + (a_{45}^{''})^{(9)}(T_{45})\right]G_{45} = 0$ | |
| $(a_{46})^{(9)}G_{45} - \left[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45})\right]G_{46} = 0$ | |
| $(b_{44})^{(9)}T_{45} - [(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}(G_{47})\,]T_{44} = 0$ | |
| $(b_{45})^{(9)}T_{44} - [(b_{45}^{'})^{(9)} - (b_{45}^{''})^{(9)}(G_{47})\,]T_{45} = 0$ | |
| $(b_{46})^{(9)}T_{45} - [(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}(G_{47})\,]T_{46} = 0$ | |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{13}, G_{14}$ if<br><br>$F(T) = (a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a_{13}^{'})^{(1)}(a_{14}^{''})^{(1)}(T_{14}) + (a_{14}^{'})^{(1)}(a_{13}^{''})^{(1)}(T_{14}) +$ | 485 |

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| $(a_{13}^{''})^{(1)}(T_{14})(a_{14}^{''})^{(1)}(T_{14}) = 0$ | |
| **Proof:**<br><br>(d) Indeed the first two equations have a nontrivial solution $G_{16}, G_{17}$ if<br><br>$\mathrm{F}(T_{19}) = (a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a_{16}^{'})^{(2)}(a_{17}^{''})^{(2)}(T_{17}) + (a_{17}^{'})^{(2)}(a_{16}^{''})^{(2)}(T_{17}) + (a_{16}^{''})^{(2)}(T_{17})(a_{17}^{''})^{(2)}(T_{17}) = 0$ | 486 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{20}, G_{21}$ if<br><br>$F(T_{23}) = (a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a_{20}^{'})^{(3)}(a_{21}^{''})^{(3)}(T_{21}) + (a_{21}^{'})^{(3)}(a_{20}^{''})^{(3)}(T_{21}) + (a_{20}^{''})^{(3)}(T_{21})(a_{21}^{''})^{(3)}(T_{21}) = 0$ | 487 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{24}, G_{25}$ if<br><br>$F(T_{27}) = (a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a_{24}^{'})^{(4)}(a_{25}^{''})^{(4)}(T_{25}) + (a_{25}^{'})^{(4)}(a_{24}^{''})^{(4)}(T_{25}) + (a_{24}^{''})^{(4)}(T_{25})(a_{25}^{''})^{(4)}(T_{25}) = 0$ | 488 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{28}, G_{29}$ if<br><br>$F(T_{31}) = (a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a_{28}^{'})^{(5)}(a_{29}^{''})^{(5)}(T_{29}) + (a_{29}^{'})^{(5)}(a_{28}^{''})^{(5)}(T_{29}) + (a_{28}^{''})^{(5)}(T_{29})(a_{29}^{''})^{(5)}(T_{29}) = 0$ | 489 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{32}, G_{33}$ if<br><br>$F(T_{35}) = (a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32}^{'})^{(6)}(a_{33}^{''})^{(6)}(T_{33}) + (a_{33}^{'})^{(6)}(a_{32}^{''})^{(6)}(T_{33}) + (a_{32}^{''})^{(6)}(T_{33})(a_{33}^{''})^{(6)}(T_{33}) = 0$ | 490 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{36}, G_{37}$ if<br><br>$F(T_{39}) = (a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36}^{'})^{(7)}(a_{37}^{''})^{(7)}(T_{37}) + (a_{37}^{'})^{(7)}(a_{36}^{''})^{(7)}(T_{37}) + (a_{36}^{''})^{(7)}(T_{37})(a_{37}^{''})^{(7)}(T_{37}) = 0$ | 491 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{40}, G_{41}$ if<br><br>$F(T_{43}) = (a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40}^{'})^{(8)}(a_{41}^{''})^{(8)}(T_{41}) + (a_{41}^{'})^{(8)}(a_{40}^{''})^{(8)}(T_{41}) + (a_{40}^{''})^{(8)}(T_{41})(a_{41}^{''})^{(8)}(T_{41}) = 0$ | 492 |

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| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{44}, G_{45}$ if<br><br>$F(T_{47}) = (a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44}^{'})^{(9)}(a_{45}^{''})^{(9)}(T_{45}) + (a_{45}^{'})^{(9)}(a_{44}^{''})^{(9)}(T_{45}) + (a_{44}^{''})^{(9)}(T_{45})(a_{45}^{''})^{(9)}(T_{45}) = 0$ | 492<br>A |
| **Definition and uniqueness of** $T_{14}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{14})$ being increasing, it follows that there exists a unique $T_{14}^{*}$ for which $f(T_{14}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{13} = \frac{(a_{13})^{(1)}G_{14}}{[(a_{13}^{'})^{(1)}+(a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15} = \frac{(a_{15})^{(1)}G_{14}}{[(a_{15}^{'})^{(1)}+(a_{15}^{''})^{(1)}(T_{14}^{*})]}$ | 493 |
| **Definition and uniqueness of** $T_{17}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(2)}(T_{17})$ being increasing, it follows that there exists a unique $T_{17}^{*}$ for which $f(T_{17}^{*}) = 0$. With this value , we obtain from the three first equations | 494 |
| $G_{16} = \frac{(a_{16})^{(2)}G_{17}}{[(a_{16}^{'})^{(2)}+(a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18} = \frac{(a_{18})^{(2)}G_{17}}{[(a_{18}^{'})^{(2)}+(a_{18}^{''})^{(2)}(T_{17}^{*})]}$ | 495 |
| **Definition and uniqueness of** $T_{21}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{21})$ being increasing, it follows that there exists a unique $T_{21}^{*}$ for which $f(T_{21}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{20} = \frac{(a_{20})^{(3)}G_{21}}{[(a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21}^{*})]}$ , $G_{22} = \frac{(a_{22})^{(3)}G_{21}}{[(a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21}^{*})]}$ | 496 |
| **Definition and uniqueness of** $T_{25}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(4)}(T_{25})$ being increasing, it follows that there exists a unique $T_{25}^{*}$ for which $f(T_{25}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{24} = \frac{(a_{24})^{(4)}G_{25}}{[(a_{24}^{'})^{(4)}+(a_{24}^{''})^{(4)}(T_{25}^{*})]}$ , $G_{26} = \frac{(a_{26})^{(4)}G_{25}}{[(a_{26}^{'})^{(4)}+(a_{26}^{''})^{(4)}(T_{25}^{*})]}$ | 497 |
| **Definition and uniqueness of** $T_{29}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(5)}(T_{29})$ being increasing, it follows that there exists a unique $T_{29}^{*}$ for which $f(T_{29}^{*}) = 0$. With this value , we obtain from the three first equations | 498 |

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| $G_{28} = \frac{(a_{28})^{(5)}G_{29}}{[(a'_{28})^{(5)}+(a''_{28})^{(5)}(T^*_{29})]}$ , $G_{30} = \frac{(a_{30})^{(5)}G_{29}}{[(a'_{30})^{(5)}+(a''_{30})^{(5)}(T^*_{29})]}$ | |
| **Definition and uniqueness of** $T^*_{33}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a''_i)^{(6)}(T_{33})$ being increasing, it follows that there exists a unique $T^*_{33}$ for which $f(T^*_{33}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{32} = \frac{(a_{32})^{(6)}G_{33}}{[(a'_{32})^{(6)}+(a''_{32})^{(6)}(T^*_{33})]}$ , $G_{34} = \frac{(a_{34})^{(6)}G_{33}}{[(a'_{34})^{(6)}+(a''_{34})^{(6)}(T^*_{33})]}$ | 499 |
| **Definition and uniqueness of** $T^*_{37}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a''_i)^{(7)}(T_{37})$ being increasing, it follows that there exists a unique $T^*_{37}$ for which $f(T^*_{37}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{36} = \frac{(a_{36})^{(7)}G_{37}}{[(a'_{36})^{(7)}+(a''_{36})^{(7)}(T^*_{37})]}$ , $G_{38} = \frac{(a_{38})^{(7)}G_{37}}{[(a'_{38})^{(7)}+(a''_{38})^{(7)}(T^*_{37})]}$ | 500 |
| **Definition and uniqueness of** $T^*_{41}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a''_i)^{(8)}(T_{41})$ being increasing, it follows that there exists a unique $T^*_{41}$ for which $f(T^*_{41}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{40} = \frac{(a_{40})^{(8)}G_{41}}{[(a'_{40})^{(8)}+(a''_{40})^{(8)}(T^*_{41})]}$ , $G_{42} = \frac{(a_{42})^{(8)}G_{41}}{[(a'_{42})^{(8)}+(a''_{42})^{(8)}(T^*_{41})]}$ | 501 |
| **Definition and uniqueness of** $T^*_{45}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a''_i)^{(9)}(T_{45})$ being increasing, it follows that there exists a unique $T^*_{45}$ for which $f(T^*_{45}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{44} = \frac{(a_{44})^{(9)}G_{45}}{[(a'_{44})^{(9)}+(a''_{44})^{(9)}(T^*_{45})]}$ , $G_{46} = \frac{(a_{46})^{(9)}G_{45}}{[(a'_{46})^{(9)}+(a''_{46})^{(9)}(T^*_{45})]}$ | 501 A |
| By the same argument, the equations admit solutions $G_{13}, G_{14}$ if<br>$\varphi(G) = (b'_{13})^{(1)}(b'_{14})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} -$<br><br>$\left[(b'_{13})^{(1)}(b''_{14})^{(1)}(G) + (b'_{14})^{(1)}(b''_{13})^{(1)}(G)\right]+(b''_{13})^{(1)}(G)(b''_{14})^{(1)}(G) = 0$<br><br>Where in $G(G_{13}, G_{14}, G_{15}), G_{13}, G_{15}$ must be replaced by their values from 96. It is easy to see that $\varphi$ is a decreasing function in $G_{14}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that | 502 |

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| there exists a unique $G_{14}^{*}$ such that $\varphi(G^{*})=0$ | |
| By the same argument, the equations admit solutions $G_{16}, G_{17}$ if<br>$\varphi(G_{19})=(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)}-(b_{16})^{(2)}(b_{17})^{(2)}-$<br>$\left[(b_{16}^{'})^{(2)}(b_{17}^{''})^{(2)}(G_{19})+(b_{17}^{'})^{(2)}(b_{16}^{''})^{(2)}(G_{19})\right]+(b_{16}^{''})^{(2)}(G_{19})(b_{17}^{''})^{(2)}(G_{19})=0$ | 503 |
| Where in $(G_{19})(G_{16},G_{17},G_{18}), G_{16}, G_{18}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{17}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{14}^{*}$ such that $\varphi((G_{19})^{*})=0$ | 504 |
| By the same argument, the equations admit solutions $G_{20}, G_{21}$ if<br>$\varphi(G_{23})=(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)}-(b_{20})^{(3)}(b_{21})^{(3)}-$<br>$\left[(b_{20}^{'})^{(3)}(b_{21}^{''})^{(3)}(G_{23})+(b_{21}^{'})^{(3)}(b_{20}^{''})^{(3)}(G_{23})\right]+(b_{20}^{''})^{(3)}(G_{23})(b_{21}^{''})^{(3)}(G_{23})=0$<br>Where in $G_{23}(G_{20},G_{21},G_{22}), G_{20}, G_{22}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{21}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{21}^{*}$ such that $\varphi((G_{23})^{*})=0$ | 505 |
| By the same argument, the equations admit solutions $G_{24}, G_{25}$ if<br>$\varphi(G_{27})=(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)}-(b_{24})^{(4)}(b_{25})^{(4)}-$<br>$\left[(b_{24}^{'})^{(4)}(b_{25}^{''})^{(4)}(G_{27})+(b_{25}^{'})^{(4)}(b_{24}^{''})^{(4)}(G_{27})\right]+(b_{24}^{''})^{(4)}(G_{27})(b_{25}^{''})^{(4)}(G_{27})=0$<br>Where in $(G_{27})(G_{24},G_{25},G_{26}), G_{24}, G_{26}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{25}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{25}^{*}$ such that $\varphi((G_{27})^{*})=0$ | 506 |
| By the same argument, the equations admit solutions $G_{28}, G_{29}$ if<br>$\varphi(G_{31})=(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)}-(b_{28})^{(5)}(b_{29})^{(5)}-$<br>$\left[(b_{28}^{'})^{(5)}(b_{29}^{''})^{(5)}(G_{31})+(b_{29}^{'})^{(5)}(b_{28}^{''})^{(5)}(G_{31})\right]+(b_{28}^{''})^{(5)}(G_{31})(b_{29}^{''})^{(5)}(G_{31})=0$<br>Where in $(G_{31})(G_{28},G_{29},G_{30}), G_{28}, G_{30}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{29}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{29}^{*}$ such that $\varphi((G_{31})^{*})=0$ | 507 |
| By the same argument, the equations admit solutions $G_{32}, G_{33}$ if<br>$\varphi(G_{35})=(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)}-(b_{32})^{(6)}(b_{33})^{(6)}-$<br>$\left[(b_{32}^{'})^{(6)}(b_{33}^{''})^{(6)}(G_{35})+(b_{33}^{'})^{(6)}(b_{32}^{''})^{(6)}(G_{35})\right]+(b_{32}^{''})^{(6)}(G_{35})(b_{33}^{''})^{(6)}(G_{35})=0$<br>Where in $(G_{35})(G_{32},G_{33},G_{34}), G_{32}, G_{34}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{33}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that | 508 |

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| there exists a unique $G_{33}^{*}$ such that $\varphi({G_{35}}^{*}) = 0$ | |
| By the same argument, the equations admit solutions $G_{36}, G_{37}$ if<br>$\varphi(G_{39}) = (b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} -$<br>$\left[(b_{36}^{'})^{(7)}(b_{37}^{''})^{(7)}(G_{39}) + (b_{37}^{'})^{(7)}(b_{36}^{''})^{(7)}(G_{39})\right] + (b_{36}^{''})^{(7)}(G_{39})(b_{37}^{''})^{(7)}(G_{39}) = 0$<br><br>Where in $(G_{39})(G_{36}, G_{37}, G_{38}), G_{36}, G_{38}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{37}$ taking into account the hypothesis $\varphi(0) > 0\,, \varphi(\infty) < 0$ it follows that there exists a unique $G_{37}^{*}$ such that $\varphi({G_{39}}^{*}) = 0$ | 509 |
| By the same argument, the equations admit solutions $G_{40}, G_{41}$ if<br>$\varphi(G_{43}) = (b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} -$<br>$\left[(b_{40}^{'})^{(8)}(b_{41}^{''})^{(8)}(G_{43}) + (b_{41}^{'})^{(8)}(b_{40}^{''})^{(8)}(G_{43})\right] + (b_{40}^{''})^{(8)}(G_{43})(b_{41}^{''})^{(8)}(G_{43}) = 0$<br><br>Where in $(G_{43})(G_{40}, G_{41}, G_{42}), G_{40}, G_{42}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{41}$ taking into account the hypothesis $\varphi(0) > 0\,, \varphi(\infty) < 0$ it follows that there exists a unique $G_{41}^{*}$ such that $\varphi({G_{43}}^{*}) = 0$ | 510 |
| By the same argument, the equations 92,93 admit solutions $G_{44}, G_{45}$ if<br>$\varphi(G_{47}) = (b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} -$<br><br>$\left[(b_{44}^{'})^{(9)}(b_{45}^{''})^{(9)}(G_{47}) + (b_{45}^{'})^{(9)}(b_{44}^{''})^{(9)}(G_{47})\right] + (b_{44}^{''})^{(9)}(G_{47})(b_{45}^{''})^{(9)}(G_{47}) = 0$<br><br>Where in $(G_{47})(G_{44}, G_{45}, G_{46}), G_{44}, G_{46}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{45}$ taking into account the hypothesis $\varphi(0) > 0\,, \varphi(\infty) < 0$ it follows that there exists a unique $G_{45}^{*}$ such that $\varphi((G_{47})^{*}) = 0$ | |
| Finally we obtain the unique solution<br><br>$G_{14}^{*}$ given by $\varphi(G^{*}) = 0\,, T_{14}^{*}$ given by $f(T_{14}^{*}) = 0$ and<br><br>$G_{13}^{*} = \frac{(a_{13})^{(1)}G_{14}^{*}}{[(a_{13}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15}^{*} = \frac{(a_{15})^{(1)}G_{14}^{*}}{[(a_{15}^{'})^{(1)} + (a_{15}^{''})^{(1)}(T_{14}^{*})]}$<br><br>$T_{13}^{*} = \frac{(b_{13})^{(1)}T_{14}^{*}}{[(b_{13}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G^{*})]}$ , $T_{15}^{*} = \frac{(b_{15})^{(1)}T_{14}^{*}}{[(b_{15}^{'})^{(1)} - (b_{15}^{''})^{(1)}(G^{*})]}$<br><br>Obviously, these values represent an equilibrium solution | 511 |
| Finally we obtain the unique solution | |
| $G_{17}^{*}$ given by $\varphi((G_{19})^{*}) = 0\,, T_{17}^{*}$ given by $f(T_{17}^{*}) = 0$ and | 512 |
| $G_{16}^{*} = \frac{(a_{16})^{(2)}G_{17}^{*}}{[(a_{16}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18}^{*} = \frac{(a_{18})^{(2)}G_{17}^{*}}{[(a_{18}^{'})^{(2)} + (a_{18}^{''})^{(2)}(T_{17}^{*})]}$ | 513 |
| $T_{16}^{*} = \frac{(b_{16})^{(2)}T_{17}^{*}}{[(b_{16}^{'})^{(2)} - (b_{16}^{''})^{(2)}((G_{19})^{*})]}$ , $T_{18}^{*} = \frac{(b_{18})^{(2)}T_{17}^{*}}{[(b_{18}^{'})^{(2)} - (b_{18}^{''})^{(2)}((G_{19})^{*})]}$ | 514 |

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| Obviously, these values represent an equilibrium solution | |
| Finally we obtain the unique solution<br>$G_{21}^*$ given by $\varphi((G_{23})^*) = 0$ , $T_{21}^*$ given by $f(T_{21}^*) = 0$ and<br>$G_{20}^* = \frac{(a_{20})^{(3)} G_{21}^*}{[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21}^*)]}$ , $G_{22}^* = \frac{(a_{22})^{(3)} G_{21}^*}{[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21}^*)]}$<br>$T_{20}^* = \frac{(b_{20})^{(3)} T_{21}^*}{[(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23}{}^*)]}$ , $T_{22}^* = \frac{(b_{22})^{(3)} T_{21}^*}{[(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23}{}^*)]}$<br>Obviously, these values represent an equilibrium solution of global equations | 515 |
| Finally we obtain the unique solution<br>$G_{25}^*$ given by $\varphi(G_{27}) = 0$ , $T_{25}^*$ given by $f(T_{25}^*) = 0$ and<br>$G_{24}^* = \frac{(a_{24})^{(4)} G_{25}^*}{[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25}^*)]}$ , $G_{26}^* = \frac{(a_{26})^{(4)} G_{25}^*}{[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25}^*)]}$ | 516 |
| $T_{24}^* = \frac{(b_{24})^{(4)} T_{25}^*}{[(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}((G_{27})^*)]}$ , $T_{26}^* = \frac{(b_{26})^{(4)} T_{25}^*}{[(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}((G_{27})^*)]}$<br>Obviously, these values represent an equilibrium solution of global equations | 517 |
| Finally we obtain the unique solution<br>$G_{29}^*$ given by $\varphi((G_{31})^*) = 0$ , $T_{29}^*$ given by $f(T_{29}^*) = 0$ and<br>$G_{28}^* = \frac{(a_{28})^{(5)} G_{29}^*}{[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}^*)]}$ , $G_{30}^* = \frac{(a_{30})^{(5)} G_{29}^*}{[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29}^*)]}$ | 518 |
| $T_{28}^* = \frac{(b_{28})^{(5)} T_{29}^*}{[(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}((G_{31})^*)]}$ , $T_{30}^* = \frac{(b_{30})^{(5)} T_{29}^*}{[(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}((G_{31})^*)]}$<br>Obviously, these values represent an equilibrium solution of global equations | 519 |
| Finally we obtain the unique solution<br>$G_{33}^*$ given by $\varphi((G_{35})^*) = 0$ , $T_{33}^*$ given by $f(T_{33}^*) = 0$ and<br>$G_{32}^* = \frac{(a_{32})^{(6)} G_{33}^*}{[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33}^*)]}$ , $G_{34}^* = \frac{(a_{34})^{(6)} G_{33}^*}{[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33}^*)]}$ | 520 |
| $T_{32}^* = \frac{(b_{32})^{(6)} T_{33}^*}{[(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}((G_{35})^*)]}$ , $T_{34}^* = \frac{(b_{34})^{(6)} T_{33}^*}{[(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}((G_{35})^*)]}$<br>Obviously, these values represent an equilibrium solution of global equations | 521 |
| Finally we obtain the unique solution<br>$G_{37}^*$ given by $\varphi((G_{39})^*) = 0$ , $T_{37}^*$ given by $f(T_{37}^*) = 0$ and | 522 |

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| $G_{36}^{*} = \frac{(a_{36})^{(7)} G_{37}^{*}}{[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)} (T_{37}^{*})]}$ , $G_{38}^{*} = \frac{(a_{38})^{(7)} G_{37}^{*}}{[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)} (T_{37}^{*})]}$<br><br>$T_{36}^{*} = \frac{(b_{36})^{(7)} T_{37}^{*}}{[(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)} ((G_{39})^{*})]}$ , $T_{38}^{*} = \frac{(b_{38})^{(7)} T_{37}^{*}}{[(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)} ((G_{39})^{*})]}$ | |
| Finally we obtain the unique solution<br><br>$G_{41}^{*}$ given by $\varphi((G_{43})^{*}) = 0$ , $T_{41}^{*}$ given by $f(T_{41}^{*}) = 0$ and<br><br>$G_{40}^{*} = \frac{(a_{40})^{(8)} G_{41}^{*}}{[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)} (T_{41}^{*})]}$ , $G_{42}^{*} = \frac{(a_{42})^{(8)} G_{41}^{*}}{[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)} (T_{41}^{*})]}$<br><br>$T_{40}^{*} = \frac{(b_{40})^{(8)} T_{41}^{*}}{[(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)} ((G_{43})^{*})]}$ , $T_{42}^{*} = \frac{(b_{42})^{(8)} T_{41}^{*}}{[(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)} ((G_{43})^{*})]}$ | 523 |
| Finally we obtain the unique solution of 89 to 99<br><br>$G_{45}^{*}$ given by $\varphi((G_{47})^{*}) = 0$ , $T_{45}^{*}$ given by $f(T_{45}^{*}) = 0$ and<br><br>$G_{44}^{*} = \frac{(a_{44})^{(9)} G_{45}^{*}}{[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)} (T_{45}^{*})]}$ , $G_{46}^{*} = \frac{(a_{46})^{(9)} G_{45}^{*}}{[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)} (T_{45}^{*})]}$<br><br>$T_{44}^{*} = \frac{(b_{44})^{(9)} T_{45}^{*}}{[(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)} ((G_{47})^{*})]}$ , $T_{46}^{*} = \frac{(b_{46})^{(9)} T_{45}^{*}}{[(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)} ((G_{47})^{*})]}$ | 523 A |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(1)}$ *and* $(b_i^{''})^{(1)}$ Belong to $C^{(1)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>**Proof:**_Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^{*} + \mathbb{G}_i$ , $T_i = T_i^{*} + \mathbb{T}_i$<br><br>$\frac{\partial (a_{14}^{''})^{(1)}}{\partial T_{14}} (T_{14}^{*}) = (q_{14})^{(1)}$ , $\frac{\partial (b_i^{''})^{(1)}}{\partial G_j} (G^{*}) = s_{ij}$ | 524 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{13}}{dt} = -\left((a_{13}^{'})^{(1)} + (p_{13})^{(1)}\right)\mathbb{G}_{13} + (a_{13})^{(1)} \mathbb{G}_{14} - (q_{13})^{(1)} G_{13}^{*} \mathbb{T}_{14}$ | 525 |

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| $\frac{d\mathbb{G}_{14}}{dt} = -\left((a_{14}^{'})^{(1)} + (p_{14})^{(1)}\right)\mathbb{G}_{14} + (a_{14})^{(1)}\mathbb{G}_{13} - (q_{14})^{(1)}G_{14}^{*}\mathbb{T}_{14}$ | 526 |
| $\frac{d\mathbb{G}_{15}}{dt} = -\left((a_{15}^{'})^{(1)} + (p_{15})^{(1)}\right)\mathbb{G}_{15} + (a_{15})^{(1)}\mathbb{G}_{14} - (q_{15})^{(1)}G_{15}^{*}\mathbb{T}_{14}$ | 527 |
| $\frac{d\mathbb{T}_{13}}{dt} = -\left((b_{13}^{'})^{(1)} - (r_{13})^{(1)}\right)\mathbb{T}_{13} + (b_{13})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(13)(j)}T_{13}^{*}\mathbb{G}_{j}\right)$ | 528 |
| $\frac{d\mathbb{T}_{14}}{dt} = -\left((b_{14}^{'})^{(1)} - (r_{14})^{(1)}\right)\mathbb{T}_{14} + (b_{14})^{(1)}\mathbb{T}_{13} + \sum_{j=13}^{15}\left(s_{(14)(j)}T_{14}^{*}\mathbb{G}_{j}\right)$ | 529 |
| $\frac{d\mathbb{T}_{15}}{dt} = -\left((b_{15}^{'})^{(1)} - (r_{15})^{(1)}\right)\mathbb{T}_{15} + (b_{15})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(15)(j)}T_{15}^{*}\mathbb{G}_{j}\right)$ | 530 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}$ Belong to $C^{(2)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable | 531 |
| Proof: Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- | |
| $G_i = G_i^{*} + \mathbb{G}_i \quad , T_i = T_i^{*} + \mathbb{T}_i$ | 532 |
| $\frac{\partial(a_{17}^{''})^{(2)}}{\partial T_{17}}(T_{17}^{*}) = (q_{17})^{(2)}$ , $\frac{\partial(b_i^{''})^{(2)}}{\partial G_j}((G_{19})^{*}) = s_{ij}$ | 533 |
| taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{16}}{dt} = -\left((a_{16}^{'})^{(2)} + (p_{16})^{(2)}\right)\mathbb{G}_{16} + (a_{16})^{(2)}\mathbb{G}_{17} - (q_{16})^{(2)}G_{16}^{*}\mathbb{T}_{17}$ | 534 |
| $\frac{d\mathbb{G}_{17}}{dt} = -\left((a_{17}^{'})^{(2)} + (p_{17})^{(2)}\right)\mathbb{G}_{17} + (a_{17})^{(2)}\mathbb{G}_{16} - (q_{17})^{(2)}G_{17}^{*}\mathbb{T}_{17}$ | 535 |
| $\frac{d\mathbb{G}_{18}}{dt} = -\left((a_{18}^{'})^{(2)} + (p_{18})^{(2)}\right)\mathbb{G}_{18} + (a_{18})^{(2)}\mathbb{G}_{17} - (q_{18})^{(2)}G_{18}^{*}\mathbb{T}_{17}$ | 536 |
| $\frac{d\mathbb{T}_{16}}{dt} = -\left((b_{16}^{'})^{(2)} - (r_{16})^{(2)}\right)\mathbb{T}_{16} + (b_{16})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(16)(j)}T_{16}^{*}\mathbb{G}_{j}\right)$ | 537 |
| $\frac{d\mathbb{T}_{17}}{dt} = -\left((b_{17}^{'})^{(2)} - (r_{17})^{(2)}\right)\mathbb{T}_{17} + (b_{17})^{(2)}\mathbb{T}_{16} + \sum_{j=16}^{18}\left(s_{(17)(j)}T_{17}^{*}\mathbb{G}_{j}\right)$ | 538 |
| $\frac{d\mathbb{T}_{18}}{dt} = -\left((b_{18}^{'})^{(2)} - (r_{18})^{(2)}\right)\mathbb{T}_{18} + (b_{18})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(18)(j)}T_{18}^{*}\mathbb{G}_{j}\right)$ | 539 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(3)}$ *and* $(b_i^{''})^{(3)}$ Belong to $C^{(3)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- | 540 |

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| $G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br>$\frac{\partial (a_{21}^{''})^{(3)}}{\partial T_{21}}(T_{21}^*) = (q_{21})^{(3)} \ , \ \frac{\partial (b_i^{''})^{(3)}}{\partial G_j}((G_{23})^*) = s_{ij}$ | |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{20}}{dt} = -\left((a_{20}^{'})^{(3)} + (p_{20})^{(3)}\right)\mathbb{G}_{20} + (a_{20})^{(3)}\mathbb{G}_{21} - (q_{20})^{(3)}G_{20}^*\mathbb{T}_{21}$ | 541 |
| $\frac{d\mathbb{G}_{21}}{dt} = -\left((a_{21}^{'})^{(3)} + (p_{21})^{(3)}\right)\mathbb{G}_{21} + (a_{21})^{(3)}\mathbb{G}_{20} - (q_{21})^{(3)}G_{21}^*\mathbb{T}_{21}$ | 542 |
| $\frac{d\mathbb{G}_{22}}{dt} = -\left((a_{22}^{'})^{(3)} + (p_{22})^{(3)}\right)\mathbb{G}_{22} + (a_{22})^{(3)}\mathbb{G}_{21} - (q_{22})^{(3)}G_{22}^*\mathbb{T}_{21}$ | 543 |
| $\frac{d\mathbb{T}_{20}}{dt} = -\left((b_{20}^{'})^{(3)} - (r_{20})^{(3)}\right)\mathbb{T}_{20} + (b_{20})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(20)(j)}T_{20}^*\mathbb{G}_j\right)$ | 544 |
| $\frac{d\mathbb{T}_{21}}{dt} = -\left((b_{21}^{'})^{(3)} - (r_{21})^{(3)}\right)\mathbb{T}_{21} + (b_{21})^{(3)}\mathbb{T}_{20} + \sum_{j=20}^{22}\left(s_{(21)(j)}T_{21}^*\mathbb{G}_j\right)$ | 545 |
| $\frac{d\mathbb{T}_{22}}{dt} = -\left((b_{22}^{'})^{(3)} - (r_{22})^{(3)}\right)\mathbb{T}_{22} + (b_{22})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(22)(j)}T_{22}^*\mathbb{G}_j\right)$ | 546 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(4)}$ *and* $(b_i^{''})^{(4)}$ Belong to $C^{(4)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 547 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br>$\frac{\partial (a_{25}^{''})^{(4)}}{\partial T_{25}}(T_{25}^*) = (q_{25})^{(4)} \ , \ \frac{\partial (b_i^{''})^{(4)}}{\partial G_j}((G_{27})^*) = s_{ij}$ | 548 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{24}}{dt} = -\left((a_{24}^{'})^{(4)} + (p_{24})^{(4)}\right)\mathbb{G}_{24} + (a_{24})^{(4)}\mathbb{G}_{25} - (q_{24})^{(4)}G_{24}^*\mathbb{T}_{25}$ | 549 |
| $\frac{d\mathbb{G}_{25}}{dt} = -\left((a_{25}^{'})^{(4)} + (p_{25})^{(4)}\right)\mathbb{G}_{25} + (a_{25})^{(4)}\mathbb{G}_{24} - (q_{25})^{(4)}G_{25}^*\mathbb{T}_{25}$ | 550 |
| $\frac{d\mathbb{G}_{26}}{dt} = -\left((a_{26}^{'})^{(4)} + (p_{26})^{(4)}\right)\mathbb{G}_{26} + (a_{26})^{(4)}\mathbb{G}_{25} - (q_{26})^{(4)}G_{26}^*\mathbb{T}_{25}$ | 551 |
| $\frac{d\mathbb{T}_{24}}{dt} = -\left((b_{24}^{'})^{(4)} - (r_{24})^{(4)}\right)\mathbb{T}_{24} + (b_{24})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(24)(j)}T_{24}^*\mathbb{G}_j\right)$ | 552 |
| $\frac{d\mathbb{T}_{25}}{dt} = -\left((b_{25}^{'})^{(4)} - (r_{25})^{(4)}\right)\mathbb{T}_{25} + (b_{25})^{(4)}\mathbb{T}_{24} + \sum_{j=24}^{26}\left(s_{(25)(j)}T_{25}^*\mathbb{G}_j\right)$ | 553 |
| $\frac{d\mathbb{T}_{26}}{dt} = -\left((b_{26}^{'})^{(4)} - (r_{26})^{(4)}\right)\mathbb{T}_{26} + (b_{26})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(26)(j)}T_{26}^*\mathbb{G}_j\right)$ | 554 |

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| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 5:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(5)}$ *and* $(b_i^{''})^{(5)}$ Belong to $C^{(5)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 555 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial (a_{29}^{''})^{(5)}}{\partial T_{29}}(T_{29}^*) = (q_{29})^{(5)}$ , $\frac{\partial (b_i^{''})^{(5)}}{\partial G_j}((G_{31})^*) = s_{ij}$ | 556 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{28}}{dt} = -\left((a_{28}^{'})^{(5)} + (p_{28})^{(5)}\right)\mathbb{G}_{28} + (a_{28})^{(5)}\mathbb{G}_{29} - (q_{28})^{(5)}G_{28}^*\mathbb{T}_{29}$ | 557 |
| $\frac{d\mathbb{G}_{29}}{dt} = -\left((a_{29}^{'})^{(5)} + (p_{29})^{(5)}\right)\mathbb{G}_{29} + (a_{29})^{(5)}\mathbb{G}_{28} - (q_{29})^{(5)}G_{29}^*\mathbb{T}_{29}$ | 558 |
| $\frac{d\mathbb{G}_{30}}{dt} = -\left((a_{30}^{'})^{(5)} + (p_{30})^{(5)}\right)\mathbb{G}_{30} + (a_{30})^{(5)}\mathbb{G}_{29} - (q_{30})^{(5)}G_{30}^*\mathbb{T}_{29}$ | 559 |
| $\frac{d\mathbb{T}_{28}}{dt} = -\left((b_{28}^{'})^{(5)} - (r_{28})^{(5)}\right)\mathbb{T}_{28} + (b_{28})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(28)(j)}T_{28}^*\mathbb{G}_j\right)$ | 560 |
| $\frac{d\mathbb{T}_{29}}{dt} = -\left((b_{29}^{'})^{(5)} - (r_{29})^{(5)}\right)\mathbb{T}_{29} + (b_{29})^{(5)}\mathbb{T}_{28} + \sum_{j=28}^{30}\left(s_{(29)(j)}T_{29}^*\mathbb{G}_j\right)$ | 561 |
| $\frac{d\mathbb{T}_{30}}{dt} = -\left((b_{30}^{'})^{(5)} - (r_{30})^{(5)}\right)\mathbb{T}_{30} + (b_{30})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(30)(j)}T_{30}^*\mathbb{G}_j\right)$ | 562 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 6:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(6)}$ *and* $(b_i^{''})^{(6)}$ Belong to $C^{(6)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 563 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial (a_{33}^{''})^{(6)}}{\partial T_{33}}(T_{33}^*) = (q_{33})^{(6)}$ , $\frac{\partial (b_i^{''})^{(6)}}{\partial G_j}((G_{35})^*) = s_{ij}$ | 564 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{32}}{dt} = -\left((a_{32}^{'})^{(6)} + (p_{32})^{(6)}\right)\mathbb{G}_{32} + (a_{32})^{(6)}\mathbb{G}_{33} - (q_{32})^{(6)}G_{32}^*\mathbb{T}_{33}$ | 565 |
| $\frac{d\mathbb{G}_{33}}{dt} = -\left((a_{33}^{'})^{(6)} + (p_{33})^{(6)}\right)\mathbb{G}_{33} + (a_{33})^{(6)}\mathbb{G}_{32} - (q_{33})^{(6)}G_{33}^*\mathbb{T}_{33}$ | 566 |

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| $\frac{d\mathbb{G}_{34}}{dt} = -\left((a'_{34})^{(6)} + (p_{34})^{(6)}\right)\mathbb{G}_{34} + (a_{34})^{(6)}\mathbb{G}_{33} - (q_{34})^{(6)}G^*_{34}\mathbb{T}_{33}$ | 567 |
| $\frac{d\mathbb{T}_{32}}{dt} = -\left((b'_{32})^{(6)} - (r_{32})^{(6)}\right)\mathbb{T}_{32} + (b_{32})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(32)(j)}T^*_{32}\mathbb{G}_j\right)$ | 568 |
| $\frac{d\mathbb{T}_{33}}{dt} = -\left((b'_{33})^{(6)} - (r_{33})^{(6)}\right)\mathbb{T}_{33} + (b_{33})^{(6)}\mathbb{T}_{32} + \sum_{j=32}^{34}\left(s_{(33)(j)}T^*_{33}\mathbb{G}_j\right)$ | 569 |
| $\frac{d\mathbb{T}_{34}}{dt} = -\left((b'_{34})^{(6)} - (r_{34})^{(6)}\right)\mathbb{T}_{34} + (b_{34})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(34)(j)}T^*_{34}\mathbb{G}_j\right)$ | 570 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 7:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(7)}$ *and* $(b_i^{''})^{(7)}$ Belong to $C^{(7)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 571 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial(a''_{37})^{(7)}}{\partial T_{37}}(T^*_{37}) = (q_{37})^{(7)}$ , $\frac{\partial(b_i^{''})^{(7)}}{\partial G_j}\left((G_{39})^{**}\right) = s_{ij}$ | 572 |
| Then taking into account equations and neglecting the terms of power 2, we obtain from | |
| $\frac{d\mathbb{G}_{36}}{dt} = -\left((a'_{36})^{(7)} + (p_{36})^{(7)}\right)\mathbb{G}_{36} + (a_{36})^{(7)}\mathbb{G}_{37} - (q_{36})^{(7)}G^*_{36}\mathbb{T}_{37}$ | 573 |
| $\frac{d\mathbb{G}_{37}}{dt} = -\left((a'_{37})^{(7)} + (p_{37})^{(7)}\right)\mathbb{G}_{37} + (a_{37})^{(7)}\mathbb{G}_{36} - (q_{37})^{(7)}G^*_{37}\mathbb{T}_{37}$ | 574 |
| $\frac{d\mathbb{G}_{38}}{dt} = -\left((a'_{38})^{(7)} + (p_{38})^{(7)}\right)\mathbb{G}_{38} + (a_{38})^{(7)}\mathbb{G}_{37} - (q_{38})^{(7)}G^*_{38}\mathbb{T}_{37}$ | 575 |
| $\frac{d\mathbb{T}_{36}}{dt} = -\left((b'_{36})^{(7)} - (r_{36})^{(7)}\right)\mathbb{T}_{36} + (b_{36})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(36)(j)}T^*_{36}\mathbb{G}_j\right)$ | 576 |
| $\frac{d\mathbb{T}_{37}}{dt} = -\left((b'_{37})^{(7)} - (r_{37})^{(7)}\right)\mathbb{T}_{37} + (b_{37})^{(7)}\mathbb{T}_{36} + \sum_{j=36}^{38}\left(s_{(37)(j)}T^*_{37}\mathbb{G}_j\right)$ | 578 |
| $\frac{d\mathbb{T}_{38}}{dt} = -\left((b'_{38})^{(7)} - (r_{38})^{(7)}\right)\mathbb{T}_{38} + (b_{38})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(38)(j)}T^*_{38}\mathbb{G}_j\right)$ | 579 |
| Obviously, these values represent an equilibrium solution<br><br>**ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 8:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(8)}$ *and* $(b_i^{''})^{(8)}$ Belong to $C^{(8)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable. | |

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| Proof: Denote | |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- $G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$ $\frac{\partial (a_{41}^{''})^{(8)}}{\partial T_{41}}(T_{41}^*) = (q_{41})^{(8)} \ , \frac{\partial (b_i^{''})^{(8)}}{\partial G_j}((G_{43})^*) = s_{ij}$ | 580 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{40}}{dt} = -\left((a_{40}^{'})^{(8)} + (p_{40})^{(8)}\right)\mathbb{G}_{40} + (a_{40})^{(8)}\mathbb{G}_{41} - (q_{40})^{(8)}G_{40}^*\mathbb{T}_{41}$ | 581 |
| $\frac{d\mathbb{G}_{41}}{dt} = -\left((a_{41}^{'})^{(8)} + (p_{41})^{(8)}\right)\mathbb{G}_{41} + (a_{41})^{(8)}\mathbb{G}_{40} - (q_{41})^{(8)}G_{41}^*\mathbb{T}_{41}$ | 582 |
| $\frac{d\mathbb{G}_{42}}{dt} = -\left((a_{42}^{'})^{(8)} + (p_{42})^{(8)}\right)\mathbb{G}_{42} + (a_{42})^{(8)}\mathbb{G}_{41} - (q_{42})^{(8)}G_{42}^*\mathbb{T}_{41}$ | 583 |
| $\frac{d\mathbb{T}_{40}}{dt} = -\left((b_{40}^{'})^{(8)} - (r_{40})^{(8)}\right)\mathbb{T}_{40} + (b_{40})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(40)(j)}T_{40}^*\mathbb{G}_j\right)$ | 584 |
| $\frac{d\mathbb{T}_{41}}{dt} = -\left((b_{41}^{'})^{(8)} - (r_{41})^{(8)}\right)\mathbb{T}_{41} + (b_{41})^{(8)}\mathbb{T}_{40} + \sum_{j=40}^{42}\left(s_{(41)(j)}T_{41}^*\mathbb{G}_j\right)$ | 585 |
| $\frac{d\mathbb{T}_{42}}{dt} = -\left((b_{42}^{'})^{(8)} - (r_{42})^{(8)}\right)\mathbb{T}_{42} + (b_{42})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(42)(j)}T_{42}^*\mathbb{G}_j\right)$ | 586 |
| **ASYMPTOTIC STABILITY ANALYSIS** **Theorem 9:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(9)}$ *and* $(b_i^{''})^{(9)}$ Belong to $C^{(9)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable. **Proof: Denote** | 586 A |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- $G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$ $\frac{\partial (a_{45}^{''})^{(9)}}{\partial T_{45}}(T_{45}^*) = (q_{45})^{(9)} \ , \frac{\partial (b_i^{''})^{(9)}}{\partial G_j}((G_{47})^*) = s_{ij}$ | |
| Then taking into account equations 89 to 99 and neglecting the terms of power 2, we obtain from 99 to 44 | |
| $\frac{d\mathbb{G}_{44}}{dt} = -\left((a_{44}^{'})^{(9)} + (p_{44})^{(9)}\right)\mathbb{G}_{44} + (a_{44})^{(9)}\mathbb{G}_{45} - (q_{44})^{(9)}G_{44}^*\mathbb{T}_{45}$ | 586 B |
| $\frac{d\mathbb{G}_{45}}{dt} = -\left((a_{45}^{'})^{(9)} + (p_{45})^{(9)}\right)\mathbb{G}_{45} + (a_{45})^{(9)}\mathbb{G}_{44} - (q_{45})^{(9)}G_{45}^*\mathbb{T}_{45}$ | 586 C |
| $\frac{d\mathbb{G}_{46}}{dt} = -\left((a_{46}^{'})^{(9)} + (p_{46})^{(9)}\right)\mathbb{G}_{46} + (a_{46})^{(9)}\mathbb{G}_{45} - (q_{46})^{(9)}G_{46}^*\mathbb{T}_{45}$ | 586 D |

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| $\frac{d\mathbb{T}_{44}}{dt} = -\left((b'_{44})^{(9)} - (r_{44})^{(9)}\right)\mathbb{T}_{44} + (b_{44})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(44)(j)}T^*_{44}\mathbb{G}_j\right)$ | 586 E |
| $\frac{d\mathbb{T}_{45}}{dt} = -\left((b'_{45})^{(9)} - (r_{45})^{(9)}\right)\mathbb{T}_{45} + (b_{45})^{(9)}\mathbb{T}_{44} + \sum_{j=44}^{46}\left(s_{(45)(j)}T^*_{45}\mathbb{G}_j\right)$ | 586 F |
| $\frac{d\mathbb{T}_{46}}{dt} = -\left((b'_{46})^{(9)} - (r_{46})^{(9)}\right)\mathbb{T}_{46} + (b_{46})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(46)(j)}T^*_{46}\mathbb{G}_j\right)$ | 586 G |
| **The characteristic equation of this system is** | 587 |
| $\left((\lambda)^{(1)} + (b'_{15})^{(1)} - (r_{15})^{(1)}\right)\{\left((\lambda)^{(1)} + (a'_{15})^{(1)} + (p_{15})^{(1)}\right)$ $\left[\left(\left((\lambda)^{(1)} + (a'_{13})^{(1)} + (p_{13})^{(1)}\right)(q_{14})^{(1)}G^*_{14} + (a_{14})^{(1)}(q_{13})^{(1)}G^*_{13}\right)\right]$ $\left(\left((\lambda)^{(1)} + (b'_{13})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(14)}T^*_{14} + (b_{14})^{(1)}s_{(13),(14)}T^*_{14}\right)$ $+\left(\left((\lambda)^{(1)} + (a'_{14})^{(1)} + (p_{14})^{(1)}\right)(q_{13})^{(1)}G^*_{13} + (a_{13})^{(1)}(q_{14})^{(1)}G^*_{14}\right)$ $\left(\left((\lambda)^{(1)} + (b'_{13})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(13)}T^*_{14} + (b_{14})^{(1)}s_{(13),(13)}T^*_{13}\right)$ $\left(\left((\lambda)^{(1)}\right)^2 + \left((a'_{13})^{(1)} + (a'_{14})^{(1)} + (p_{13})^{(1)} + (p_{14})^{(1)}\right)(\lambda)^{(1)}\right)$ $\left(\left((\lambda)^{(1)}\right)^2 + \left((b'_{13})^{(1)} + (b'_{14})^{(1)} - (r_{13})^{(1)} + (r_{14})^{(1)}\right)(\lambda)^{(1)}\right)$ $+\left(\left((\lambda)^{(1)}\right)^2 + \left((a'_{13})^{(1)} + (a'_{14})^{(1)} + (p_{13})^{(1)} + (p_{14})^{(1)}\right)(\lambda)^{(1)}\right)(q_{15})^{(1)}G_{15}$ $+\left((\lambda)^{(1)} + (a'_{13})^{(1)} + (p_{13})^{(1)}\right)\left((a_{15})^{(1)}(q_{14})^{(1)}G^*_{14} + (a_{14})^{(1)}(a_{15})^{(1)}(q_{13})^{(1)}G^*_{13}\right)$ $\left(\left((\lambda)^{(1)} + (b'_{13})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(15)}T^*_{14} + (b_{14})^{(1)}s_{(13),(15)}T^*_{13}\right)\} = 0$ $+$ $\left((\lambda)^{(2)} + (b'_{18})^{(2)} - (r_{18})^{(2)}\right)\{\left((\lambda)^{(2)} + (a'_{18})^{(2)} + (p_{18})^{(2)}\right)$ $\left[\left(\left((\lambda)^{(2)} + (a'_{16})^{(2)} + (p_{16})^{(2)}\right)(q_{17})^{(2)}\mathrm{G}^*_{17} + (a_{17})^{(2)}(q_{16})^{(2)}\mathrm{G}^*_{16}\right)\right]$ $\left(\left((\lambda)^{(2)} + (b'_{16})^{(2)} - (r_{16})^{(2)}\right)s_{(17),(17)}\mathrm{T}^*_{17} + (b_{17})^{(2)}s_{(16),(17)}\mathrm{T}^*_{17}\right)$ $+\left(\left((\lambda)^{(2)} + (a'_{17})^{(2)} + (p_{17})^{(2)}\right)(q_{16})^{(2)}\mathrm{G}^*_{16} + (a_{16})^{(2)}(q_{17})^{(2)}\mathrm{G}^*_{17}\right)$ $\left(\left((\lambda)^{(2)} + (b'_{16})^{(2)} - (r_{16})^{(2)}\right)s_{(17),(16)}\mathrm{T}^*_{17} + (b_{17})^{(2)}s_{(16),(16)}\mathrm{T}^*_{16}\right)$ $\left(\left((\lambda)^{(2)}\right)^2 + \left((a'_{16})^{(2)} + (a'_{17})^{(2)} + (p_{16})^{(2)} + (p_{17})^{(2)}\right)(\lambda)^{(2)}\right)$ $\left(\left((\lambda)^{(2)}\right)^2 + \left((b'_{16})^{(2)} + (b'_{17})^{(2)} - (r_{16})^{(2)} + (r_{17})^{(2)}\right)(\lambda)^{(2)}\right)$ | |

$$+\left(\left((\lambda)^{(2)}\right)^2+\left((a_{16}')^{(2)}+(a_{17}')^{(2)}+(p_{16})^{(2)}+(p_{17})^{(2)}\right)(\lambda)^{(2)}\right)(q_{18})^{(2)}\mathrm{G}_{18}$$

$$+\left((\lambda)^{(2)}+(a_{16}')^{(2)}+(p_{16})^{(2)}\right)\left((a_{18})^{(2)}(q_{17})^{(2)}\mathrm{G}_{17}^{*}+(a_{17})^{(2)}(a_{18})^{(2)}(q_{16})^{(2)}\mathrm{G}_{16}^{*}\right)$$

$$\left(\left((\lambda)^{(2)}+(b_{16}')^{(2)}-(r_{16})^{(2)}\right)s_{(17),(18)}\mathrm{T}_{17}^{*}+(b_{17})^{(2)}s_{(16),(18)}\mathrm{T}_{16}^{*}\right)\}=0$$

$$+$$

$$\left((\lambda)^{(3)}+(b_{22}')^{(3)}-(r_{22})^{(3)}\right)\{\left((\lambda)^{(3)}+(a_{22}')^{(3)}+(p_{22})^{(3)}\right)$$

$$\left[\left(\left((\lambda)^{(3)}+(a_{20}')^{(3)}+(p_{20})^{(3)}\right)(q_{21})^{(3)}G_{21}^{*}+(a_{21})^{(3)}(q_{20})^{(3)}G_{20}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(21)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(21)}T_{21}^{*}\right)$$

$$+\left(\left((\lambda)^{(3)}+(a_{21}')^{(3)}+(p_{21})^{(3)}\right)(q_{20})^{(3)}G_{20}^{*}+(a_{20})^{(3)}(q_{21})^{(1)}G_{21}^{*}\right)$$

$$\left(\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(20)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(20)}T_{20}^{*}\right)$$

$$\left(\left((\lambda)^{(3)}\right)^2+\left((a_{20}')^{(3)}+(a_{21}')^{(3)}+(p_{20})^{(3)}+(p_{21})^{(3)}\right)(\lambda)^{(3)}\right)$$

$$\left(\left((\lambda)^{(3)}\right)^2+\left((b_{20}')^{(3)}+(b_{21}')^{(3)}-(r_{20})^{(3)}+(r_{21})^{(3)}\right)(\lambda)^{(3)}\right)$$

$$+\left(\left((\lambda)^{(3)}\right)^2+\left((a_{20}')^{(3)}+(a_{21}')^{(3)}+(p_{20})^{(3)}+(p_{21})^{(3)}\right)(\lambda)^{(3)}\right)(q_{22})^{(3)}G_{22}$$

$$+\left((\lambda)^{(3)}+(a_{20}')^{(3)}+(p_{20})^{(3)}\right)\left((a_{22})^{(3)}(q_{21})^{(3)}G_{21}^{*}+(a_{21})^{(3)}(a_{22})^{(3)}(q_{20})^{(3)}G_{20}^{*}\right)$$

$$\left(\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(22)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(22)}T_{20}^{*}\right)\}=0$$

$$+$$

$$\left((\lambda)^{(4)}+(b_{26}')^{(4)}-(r_{26})^{(4)}\right)\{\left((\lambda)^{(4)}+(a_{26}')^{(4)}+(p_{26})^{(4)}\right)$$

$$\left[\left(\left((\lambda)^{(4)}+(a_{24}')^{(4)}+(p_{24})^{(4)}\right)(q_{25})^{(4)}G_{25}^{*}+(a_{25})^{(4)}(q_{24})^{(4)}G_{24}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(4)}+(b_{24}')^{(4)}-(r_{24})^{(4)}\right)s_{(25),(25)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(25)}T_{25}^{*}\right)$$

$$+\left(\left((\lambda)^{(4)}+(a_{25}')^{(4)}+(p_{25})^{(4)}\right)(q_{24})^{(4)}G_{24}^{*}+(a_{24})^{(4)}(q_{25})^{(4)}G_{25}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}+(b_{24}')^{(4)}-(r_{24})^{(4)}\right)s_{(25),(24)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(24)}T_{24}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}\right)^2+\left((a_{24}')^{(4)}+(a_{25}')^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)$$

$$\left(\left((\lambda)^{(4)}\right)^2+\left((b_{24}')^{(4)}+(b_{25}')^{(4)}-(r_{24})^{(4)}+(r_{25})^{(4)}\right)(\lambda)^{(4)}\right)$$

$$+\left(\left((\lambda)^{(4)}\right)^2+\left((a_{24}^{'})^{(4)}+(a_{25}^{'})^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)(q_{26})^{(4)}G_{26}$$

$$+\left((\lambda)^{(4)}+(a_{24}^{'})^{(4)}+(p_{24})^{(4)}\right)\left((a_{26})^{(4)}(q_{25})^{(4)}G_{25}^{*}+(a_{25})^{(4)}(a_{26})^{(4)}(q_{24})^{(4)}G_{24}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}+(b_{24}^{'})^{(4)}-(r_{24})^{(4)}\right)s_{(25),(26)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(26)}T_{24}^{*}\right)\}=0$$

$$+$$

$$(\lambda)^{(5)}+(b_{30}^{'})^{(5)}-(r_{30})^{(5)}\}\{\left((\lambda)^{(5)}+(a_{30}^{'})^{(5)}+(p_{30})^{(5)}\right)$$

$$\left[\left(\left((\lambda)^{(5)}+(a_{28}^{'})^{(5)}+(p_{28})^{(5)}\right)(q_{29})^{(5)}G_{29}^{*}+(a_{29})^{(5)}(q_{28})^{(5)}G_{28}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(5)}+(b_{28}^{'})^{(5)}-(r_{28})^{(5)}\right)s_{(29),(29)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(29)}T_{29}^{*}\right)$$

$$+\left(\left((\lambda)^{(5)}+(a_{29}^{'})^{(5)}+(p_{29})^{(5)}\right)(q_{28})^{(5)}G_{28}^{*}+(a_{28})^{(5)}(q_{29})^{(5)}G_{29}^{*}\right)$$

$$\left(\left((\lambda)^{(5)}+(b_{28}^{'})^{(5)}-(r_{28})^{(5)}\right)s_{(29),(28)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(28)}T_{28}^{*}\right)$$

$$\left(\left((\lambda)^{(5)}\right)^2+\left((a_{28}^{'})^{(5)}+(a_{29}^{'})^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)$$

$$\left(\left((\lambda)^{(5)}\right)^2+\left((b_{28}^{'})^{(5)}+(b_{29}^{'})^{(5)}-(r_{28})^{(5)}+(r_{29})^{(5)}\right)(\lambda)^{(5)}\right)$$

$$+\left(\left((\lambda)^{(5)}\right)^2+\left((a_{28}^{'})^{(5)}+(a_{29}^{'})^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)(q_{30})^{(5)}G_{30}$$

$$+\left((\lambda)^{(5)}+(a_{28}^{'})^{(5)}+(p_{28})^{(5)}\right)\left((a_{30})^{(5)}(q_{29})^{(5)}G_{29}^{*}+(a_{29})^{(5)}(a_{30})^{(5)}(q_{28})^{(5)}G_{28}^{*}\right)$$

$$\left(\left((\lambda)^{(5)}+(b_{28}^{'})^{(5)}-(r_{28})^{(5)}\right)s_{(29),(30)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(30)}T_{28}^{*}\right)\}=0$$

$$+$$

$$(\lambda)^{(6)}+(b_{34}^{'})^{(6)}-(r_{34})^{(6)}\}\{\left((\lambda)^{(6)}+(a_{34}^{'})^{(6)}+(p_{34})^{(6)}\right)$$

$$\left[\left(\left((\lambda)^{(6)}+(a_{32}^{'})^{(6)}+(p_{32})^{(6)}\right)(q_{33})^{(6)}G_{33}^{*}+(a_{33})^{(6)}(q_{32})^{(6)}G_{32}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(6)}+(b_{32}^{'})^{(6)}-(r_{32})^{(6)}\right)s_{(33),(33)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(33)}T_{33}^{*}\right)$$

$$+\left(\left((\lambda)^{(6)}+(a_{33}^{'})^{(6)}+(p_{33})^{(6)}\right)(q_{32})^{(6)}G_{32}^{*}+(a_{32})^{(6)}(q_{33})^{(6)}G_{33}^{*}\right)$$

$$\left(\left((\lambda)^{(6)}+(b_{32}^{'})^{(6)}-(r_{32})^{(6)}\right)s_{(33),(32)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(32)}T_{32}^{*}\right)$$

$$\left(\left((\lambda)^{(6)}\right)^2+\left((a_{32}^{'})^{(6)}+(a_{33}^{'})^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\right)$$

$$\left(\left((\lambda)^{(6)}\right)^2+\left((b_{32}^{'})^{(6)}+(b_{33}^{'})^{(6)}-(r_{32})^{(6)}+(r_{33})^{(6)}\right)(\lambda)^{(6)}\right)$$

$$+\left(\left((\lambda)^{(6)}\right)^2+\left((a'_{32})^{(6)}+(a'_{33})^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\right)(q_{34})^{(6)}G_{34}$$

$$+\left((\lambda)^{(6)}+(a'_{32})^{(6)}+(p_{32})^{(6)}\right)\left((a_{34})^{(6)}(q_{33})^{(6)}G^*_{33}+(a_{33})^{(6)}(a_{34})^{(6)}(q_{32})^{(6)}G^*_{32}\right)$$

$$\left((\lambda)^{(6)}+(b'_{32})^{(6)}-(r_{32})^{(6)}\right)s_{(33),(34)}T^*_{33}+(b_{33})^{(6)}s_{(32),(34)}T^*_{32}\Big)\}=0$$

$$+$$

$$(\lambda)^{(7)}+(b'_{38})^{(7)}-(r_{38})^{(7)}\}\{\left((\lambda)^{(7)}+(a'_{38})^{(7)}+(p_{38})^{(7)}\right)$$

$$\left[\left(\left((\lambda)^{(7)}+(a'_{36})^{(7)}+(p_{36})^{(7)}\right)(q_{37})^{(7)}G^*_{37}+(a_{37})^{(7)}(q_{36})^{(7)}G^*_{36}\right)\right]$$

$$\left((\lambda)^{(7)}+(b'_{36})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(37)}T^*_{37}+(b_{37})^{(7)}s_{(36),(37)}T^*_{37}\Big)$$

$$+\left(\left((\lambda)^{(7)}+(a'_{37})^{(7)}+(p_{37})^{(7)}\right)(q_{36})^{(7)}G^*_{36}+(a_{36})^{(7)}(q_{37})^{(7)}G^*_{37}\right)$$

$$\left((\lambda)^{(7)}+(b'_{36})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(36)}T^*_{37}+(b_{37})^{(7)}s_{(36),(36)}T^*_{36}\Big)$$

$$\left((\lambda)^{(7)}\right)^2+\left((a'_{36})^{(7)}+(a'_{37})^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\Big)$$

$$\left((\lambda)^{(7)}\right)^2+\left((b'_{36})^{(7)}+(b'_{37})^{(7)}-(r_{36})^{(7)}+(r_{37})^{(7)}\right)(\lambda)^{(7)}\Big)$$

$$+\left(\left((\lambda)^{(7)}\right)^2+\left((a'_{36})^{(7)}+(a'_{37})^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\right)(q_{38})^{(7)}G_{38}$$

$$+\left((\lambda)^{(7)}+(a'_{36})^{(7)}+(p_{36})^{(7)}\right)\left((a_{38})^{(7)}(q_{37})^{(7)}G^*_{37}+(a_{37})^{(7)}(a_{38})^{(7)}(q_{36})^{(7)}G^*_{36}\right)$$

$$\left((\lambda)^{(7)}+(b'_{36})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(38)}T^*_{37}+(b_{37})^{(7)}s_{(36),(38)}T^*_{36}\Big)\}=0$$

$$+$$

$$(\lambda)^{(8)}+(b'_{42})^{(8)}-(r_{42})^{(8)}\}\{\left((\lambda)^{(8)}+(a'_{42})^{(8)}+(p_{42})^{(8)}\right)$$

$$\left[\left(\left((\lambda)^{(8)}+(a'_{40})^{(8)}+(p_{40})^{(8)}\right)(q_{41})^{(8)}G^*_{41}+(a_{41})^{(8)}(q_{40})^{(8)}G^*_{40}\right)\right]$$

$$\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(41)}T^*_{41}+(b_{41})^{(8)}s_{(40),(41)}T^*_{41}\Big)$$

$$+\left(\left((\lambda)^{(8)}+(a'_{41})^{(8)}+(p_{41})^{(8)}\right)(q_{40})^{(8)}G^*_{40}+(a_{40})^{(8)}(q_{41})^{(8)}G^*_{41}\right)$$

$$\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(40)}T^*_{41}+(b_{41})^{(8)}s_{(40),(40)}T^*_{40}\Big)$$

$$\left((\lambda)^{(8)}\right)^2+\left((a'_{40})^{(8)}+(a'_{41})^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\Big)$$

$$\left((\lambda)^{(8)}\right)^2+\left((b'_{40})^{(8)}+(b'_{41})^{(8)}-(r_{40})^{(8)}+(r_{41})^{(8)}\right)(\lambda)^{(8)}\Big)$$

$$+\left(\left((\lambda)^{(8)}\right)^2+\left((a'_{40})^{(8)}+(a'_{41})^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\right)(q_{42})^{(8)}G_{42}$$

$$+\left((\lambda)^{(8)}+(a'_{40})^{(8)}+(p_{40})^{(8)}\right)\left((a_{42})^{(8)}(q_{41})^{(8)}G^*_{41}+(a_{41})^{(8)}(a_{42})^{(8)}(q_{40})^{(8)}G^*_{40}\right)$$

$$\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(42)}T^*_{41}+(b_{41})^{(8)}s_{(40),(42)}T^*_{40}\Big)\}=0$$

$$+$$

$$(\lambda)^{(9)}+(b'_{46})^{(9)}-(r_{46})^{(9)}\}\{\left((\lambda)^{(9)}+(a'_{46})^{(9)}+(p_{46})^{(9)}\right)$$

$$\left[\left(\left((\lambda)^{(9)}+(a'_{44})^{(9)}+(p_{44})^{(9)}\right)(q_{45})^{(9)}G^*_{45}+(a_{45})^{(9)}(q_{44})^{(9)}G^*_{44}\right)\right]$$

$$\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(45)}T^*_{45}+(b_{45})^{(9)}s_{(44),(45)}T^*_{45}\Big)$$

$$+\left(\left((\lambda)^{(9)}+(a'_{45})^{(9)}+(p_{45})^{(9)}\right)(q_{44})^{(9)}G^*_{44}+(a_{44})^{(9)}(q_{45})^{(9)}G^*_{45}\right)$$

$$\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(44)}T^*_{45}+(b_{45})^{(9)}s_{(44),(44)}T^*_{44}\Big)$$

$$\left((\lambda)^{(9)}\right)^2+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\Big)$$

$$\left((\lambda)^{(9)}\right)^2+\left((b'_{44})^{(9)}+(b'_{45})^{(9)}-(r_{44})^{(9)}+(r_{45})^{(9)}\right)(\lambda)^{(9)}\Big)$$

$$+\left(\left((\lambda)^{(9)}\right)^2+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\right)(q_{46})^{(9)}G_{46}$$

$$+\left((\lambda)^{(9)}+(a'_{44})^{(9)}+(p_{44})^{(9)}\right)\left((a_{46})^{(9)}(q_{45})^{(9)}G^*_{45}+(a_{45})^{(9)}(a_{46})^{(9)}(q_{44})^{(9)}G^*_{44}\right)$$

$$\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(46)}T^*_{45}+(b_{45})^{(9)}s_{(44),(46)}T^*_{44}\Big)\}=0$$

**And as one sees, all the coefficients are positive. It follows that all the roots have negative real part, and this proves the theorem.**

Note: **Same dovetailing explanation holds good for d/dt (partial differential with respect to t), d/dt,d^2/dt^2 (acceralation: double dot). Such an exposition is helpful in optimisation problems, duality of motion, wave motion a homogeneous functions studies, and control theory which we intend to incorporate in future.** It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the

concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

# Persistence, Coexistence And Thrusty Permanence For Interacting Structured Particles, Transactions In Finance, Econometric Productivity Functions, Environment, Interfaces, Cosmological Interactions Et Al

## Josef Hofbauer And Sebastian J. Schreiber Model Thereof: A Functional Appraisal

ABSTRACT. **JOSEF HOFBAUER AND SEBASTIAN J. SCHREIBER** studied the dynamics of interacting structured systems (in their case population dynamics) can be modeled by $\frac{dxi}{dt} = A_i\ (x)\ xi$ where $x_i\ \epsilon\ R^n_{i,}\ x = (x_1, \ldots, x_k)$, and $A_i(x)$ are matrices with non-negative off- diagonal entries. These models are permanent if there exists a positive global attractor and are robustly permanent if they remain permanent following perturbations of $A_i\ (x)$. Necessary and sufficient conditions for robust permanence are derived using dominant Lyapunov exponents $X_i\ (\mu)$ of the $A_i\ (x)$ with respect to invariant measures $\mu$. The necessary conditions requires $\max_i\ X_i(\mu)>0$ for all ergodic measures with support in the boundary of the non-negative cone. The sufficient condition requires that the boundary admits a Morse decomposition such that $\max_i$ Xi$(\mu)>0$ for all invariant measures $\mu$ supported by a component of the Morse decomposition. When the Morse components are Axiom A, uniquely ergodic, or support all but one population, the necessary and sufficient conditions are equivalent. Applications to spatial ecology, epidemiology, and gene networks are given. Appeared in *Journal of Differential Equations*, 248, 1955-1971 (2010) It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## INTRODUCTION

A fundamental issue in population biology is what are the minimal conditions to ensure the long-term survivorship for interacting populations whether they be viral particles, bio-chemicals, plants, or animals. When these conditions are met the interacting populations are said to persist or coexist. Since the pioneering work of Lotka and Volterra on competitive and predator prey interactions, Thompson, Nicholson, and Bailey on host-parasite interactions, and

Kermack and McKendrick on disease outbreaks, nonlinear difference and differential equations have been used to understand conditions for population persistence [1-5]. One particularly important form of persistence for deterministic models is *permanence* or *uniform persistence* which corresponds to the existence of a global attractor bounded away from extinction of one or more species. When such an attractor exists, interacting populations are able to recover from "vigorous shake ups" of the of the population state. Permanence has been characterized from a topological perspective [6-8] and with average Lyapunov functions [9-12] for differential equation models of the form

$$(4) \qquad \frac{dxi}{dt} = x_i f_i(x) \qquad i = 1, \ldots, k$$

Where $x_i$ is the abundance of population $i$, $x = (x_1, \ldots, x_k)$, and $f_i(x)$ is the per-capita growth of population $i$.

An important extension of the concept of the permanence is *robust permanence* [13-16] which requires that (1) remain permanent after sufficiently small perturbations of the per-capita growth rates $f_i$. The importance of robust permanence as one crude form of structural stability for population models are approximations to reality. Consequently, if nearby models (e.g. more realistic models) are not permanent despite the focal model being permanent, and then one can draw few conclusions about the persistence of the biological system being approximated by the model. One can view robust permanence as one crude form of structural stability for debits-credits particles, propositions, transactions in the celestial subterranean realm and ceratoid dualism, finance like debiting say a current account and crediting to another account, or recrediting to the account which had been earlier debited, further conduction of the cash credit account, black hole formation, guzzling of inventories for the output, creation of income from output, reinvestment of income to investment,gormadizing of matter by black holes, and emission of radiation therefore, gamma ray bursts, in fact most of the transactions and interactions in the universe are conservative and quantum mechanical universe, something like a account acting as predator and prey. More often than not as we have seen from Haimovici model such models could be applied to the problems, propositions, which are not subsumed by the concomitant problems in the bearing. the population models. For this perspective, it is not unexpected that there are permanent systems that cannot be approximated by robustly permanent systems [17]. In [14, 16], criteria for robust permanence were developed with respect to the average per-capita growth rates $\int f_i(x)d\mu(x)$ with respect to invariant probability measures $\mu$ supported on the boundary of the positive orthant (i.e.

where one or more Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates are extinct). Roughly, these criteria for robust permanence require that the average per-capita growth rate (i.e. invasion rate) is positive for some missing species.

While (1) can account for many types of Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates. Interactions, it assumes that all individual within populations are exactly the same. However, theoretical biologists have long recognized that different individuals within a population may be in different states (e.g. different sizes or ages, living in different spatial locations) and these differences can have important consequences for population dynamics [18-24]. To account for how these differences influence persistence of interacting populations, we develop criteria for robust permanence for the dynamics of $k$ interacting structured populations. Roughly, these dynamics correspond to replacing $x_i$ in (1) by vectors and $f_i(x)$ by matrices. In section 2, we describe these models in greater detail and review some basic concepts from dynamical systems theory. In section 3, we introduce the structured analog of average per-capita growth rates and define the concept of robustly unsaturated invariant sets. Necessary and sufficient conditions for an invariant set to be robustly unsaturated are proven. For Axiom A or uniquely ergodic invariant sets, the necessary and sufficient criteria are shown to be equivalent. In section 4, we use Morse decompositions and the criteria for robustly unsaturated invariant sets to develop necessary and sufficient conditions for robust permanence. In sections 5 and 6, we provide several applications of our results to spatially structured ecological models, structured epidemiological models, and models of gene networks. We apply the model to econometrics. To give an example investment needs inventory to be used towards the consummation of output, part of output is to be sold for consolidation of income, and some income might be replenished in the investment and so on. In finance a cash credit account could be debited or credited when there is say purchase of demand bills or discount of usance bills. Cash credit account is debited whenever cheques are issued towards the purchase of raw materials and so on. In fact some of the models have been applied to Organic psychiatry and we have shown how such models are useful in dealing with the propositions of Husserl, Derrida, Deleuze and others.

## 4. MODELS AND ASSUMPTIONS

Let $x_i$ denote the state of the $i$-th population that lies in the non-negative cone $C_i$ of $R^{ni.}$ Define $x=(x_1,\ldots x_k)\epsilon C$ to be the non-negative cone of $R^n$ where $n=\sum_{i=1}^{k} n_i$. Let $C_i^+$ denote positive cone { $x_i \epsilon C_i$: $\prod_i$i noitalupop rof$^j$

$i$ *and* $C^+=\prod_i C^+_i$ , the positive cone for the interacting populations. If $x \mapsto A_i(x)$ is a map into $n_i x n_i$ matrices that describes the growth of population I, then the dynamics of the interacting populations are given by

$$(2)\frac{dx_i}{dt} = A_i(x)x_i \quad i=1,\ldots,k$$

Let $x.t$ denote the solution of (2), with the initial condition x, and more generally the semi flow generated by (2). Before stating our assumption on (2), recall a few definitions from dynamical systems. Given sets $I \subseteq R$ and $K \subseteq C$, let $K.I = \{x.t : t \epsilon I, x \epsilon K\}$. A set $K \subseteq C$ *is invariant* if K.t = K for all t>0, and *forward invariant* if $K.t \subset K$ for all $t > 0$. The *omega limit* set of a set $K \subseteq C$ equals $\omega(K) = \bigcap_{t\geq 0,} \overline{K.[t,\infty)}$. The *alpha limit* set of $K \subseteq C$ is $a(K) = \bigcap_{t\leq 0} \overline{K.(-\infty,t]}$. Given a forward invariant set K, $B \subset K$ is an attractor for the semi flow $x.t$ *restricted* to K provided there exists an open neighborhood $U \subseteq K$ of B such that $\omega(U) = B$. The stable set of a *compact invariant* set $K$ is defined by

$$W^s(K) = \{x \epsilon C : \omega(x) \neq \emptyset \text{ and } \omega(x) \subseteq K\}.$$

The semi flow generated by (2) is *dissipative* if there exists an attractor B with $W^s(B) = C$.

Throughout this paper, we make the following assumptions:

**A1**: $x \mapsto A_i(x)$ is continuous

**A2:** $A_i(x)$ is irreducible and has non-negative off- diagonal entries,

**A3:** $x.t$ is defined for all $t \geq 0$

**A4:** $x.t$ is dissipative with global attractor $\Gamma(A)$.

A1 is a basic regularity assumption. The non-negative of off- diagonal entries in A2 implies that there are no negative feedbacks between individuals of different states in population *i*. This assumption is meet for many types of structured models as discussed in sections 5 and 6. The irreducibility assumption A2 is generally meet and implies that all individuals within a population can pass through all states. In Remark1, we discuss how this irreducibility assumption can be relaxed. A3 ensures that the population dynamics are defined for all future time. A4 requires that

population densities/ abundances eventually are uniformly bounded, a condition that should be meet for an biologically realistic model.

We say that the semi flow of (2) is *permanent* if this semi flow is dissipative and there is a positive attractor $B \subset C^+$ such that $W^s(B) \subset C^+$. We will denote this positive attractor by $\Lambda(A)$. Permanence ensures that populations can recover from rare large perturbations and allows for a diversity of dynamical behaviors.

For K ⊂ C and $\delta > 0$, define the $\delta$- *neighborhood* of K as

$$N\delta(K) = \{x \in C : \backslash x - y < \delta \text{ for some } y \in K\}.$$

We define a $\delta$- *perturbation of (2)* to be a system of the form

$$\frac{dx}{dt} = \tilde{A}(x)x$$

that satisfies assumptions A1-A4, $\Gamma(\tilde{A}) \subset N$ *( $\Gamma(A)$), and* $||A(x) - \tilde{A}(x)|| \leq \delta$ *for all* $x \in N$ *($\Gamma(A)$).* (2) is *robustly permanent* if there exist >0 and $\epsilon > 0$ such that all $\delta$- perturbations of (2) are permanent, and $d(\Lambda(\tilde{A}), C\backslash C^+) > \epsilon$ for all $\delta$- perturbations i.e., there is a common/ uniform region of repulsion around the boundary.

## INUNDATION OF COMPACT SETS

We begin by studying the linear skew product flows on $\Gamma(A) \times R^n_i$ defined by $(x,y).t = (x.t, B_i(t,x)y)$ where $Y(t) = B_i(t,x)$ is the solution to $\frac{dY}{dt}(t) = A_i(x.t)Y(t)$ with y(0) equal to the identify matrix. Our assumption that $A_i$ is irreducible with non- negative off diagonal entries implies that $B_i$ (t,x)$C_i \subset C_i^+$ for all $x$ and t > 0 (see, e.g., [25] ). A result of Ruelle [26, Prop, 3.2] implies that there exist continuous maps $u_i\, v_i : \Gamma(A) \rightarrow C_i^+$ with $u_i(x)| = |v_i(x)| = 1$ such that The line bundle $E_i(x)$ spanned by $u_i(x)$ is invariant i.e. $E_i(x.t) = B_i(x,t)E_i(x)$ for all t ≥ 0.The vector bundle $F_i(x)$ perpendicular to $v_i(x)$ is invariant i.e. $F_i(x.t) = B_i(x,t)F_i(x)$ for all $t \geq 0$.There exist constants $a > 0$ *and* $\beta > 0$ such that(3) $||B_i(t,x)F_i(x)|| \leq a\, exp(-\beta t)||B_i(t,x)E_i(x)||$ For all $x \in K$ *and* $t \geq 0$.**Interaction rates.** Given any x $\epsilon$ $\Gamma$(A), define the *interaction and inundation rate* of *species i at particle/proposition/problem/ state x as* $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right)$ *in* $||B_i(t,x)||$ when $x_i = 0$, $\lambda_i(x)$ provides an upper bound to the rate of growth of population $i$ when introduced at infinitesimally small densities. Important properties of this invasion rate are summarized in the following two propositions. For instance, the second proposition implies if $(x.t)_i$ stays bounded away from zero (i.e. population $i$ persist), then

$\lambda_i(x)=0$. Alternatively, if $\lambda_i(x)<0$,then population $i$ is doomed to extinction. **Proposition 1.** $\lambda_i(x)$ *satisfies the following properties:* $\lambda_i(x)= limsup_{t\to\infty}\left(\frac{1}{t}\right) In \ |B_i \ (t,x)v|$ *for any* $v>0$ *in* $R^{ni}$, *and* $\lambda_i(x)= limsup_{t\to\infty}\left(\frac{1}{t}\right) \int^{t}_{0}(A_i(x.s)u_i(x.s),u_i(x.s))ds.$

**Proof:** To prove the first property, we first show that

$$\lambda i(x) = \underset{t\to\infty}{limsup}\frac{1}{t} In\ |B_i(t,x)u_i(x)|.$$

To this end, let $v \epsilon R^{ni}$ be any non-zero vector. Since $R^{ni}= E_{i(x)} \ F_i(x)$, we can write $v=au_i(x)+w$ with $a \epsilon R$ and $w \epsilon F_i(x)$. Equation (3) implies

$|B_i(t,x)v| \ \leq \ a|\ B_i(t,x)u_i(x)|+ |B_i(t,x)w| \ \leq \ |\ B_i(t,x)u_i(x)|(a+ \alpha+ exp\ (-\beta t)|\ w|)$. Hence, $\underset{t\to\infty}{limsup}\frac{1}{t} In\ B_i(t,x)\ v| \leq$

$$\underset{t\to\infty}{limsup}\frac{1}{t} In\ B_i\ (t,x)\ u_i\ (x) \leq \lambda_i(x)$$

for all non zero vectors $v \epsilon R^{ni}$. Since $||B_i(t,x)|| = sup|v|=1B_i(t,x)v|$, this inequality implies that $\lambda i(x) = \underset{t\to\infty}{\text{limsup}}\frac{1}{t}$ In $|B_i(t,x)v|$, this inequality implies that $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right)$ In $|t,x)u_i(x)|$. Now let v $\epsilon$ $C_i^{+}$. Then, we can write $v = au_i(x)+ w$ with $a >0$ and $w \epsilon F_i$.

Equation (3) implies

$$|B_i(t,x)v| \ \geq \ a|B_i(t,x)u_i(x)|-|B_i(t,x)w|$$

$$\geq |\ B_i(t,x)u_i(x)|(a- \alpha+ exp\ (-\beta t)|\ w|).$$

Since a>0, this inequality implies that

$$\lambda i(x) \geq \underset{t\to\infty}{\text{limsup}}\frac{1}{t} \text{In}\ |B_i(t,x)v| \ \geq \ \lambda i(x) \geq \underset{t\to\infty}{\text{limsup}}\frac{1}{t} \text{In}\ |B_i(t,x)u_i(x)| = \lambda i(x)$$

To prove the second assertion, let $b_i(t,x) = In\ |B_i(t,x)u_i(x)|$. Invariance of $u_i(x)$ implies that $b_i(t,x)$ is additive:

$$B_i(t+s,x) \ = \ In|B_i(t+s,x)\ u_i(x)|= In\ |B_i(s,x.t)B_i(t,x)u_i(x)$$

$$= \text{In}\ |B_i(s,x.t)\ u_i(x.t)||\ B_i(t,x)\ u_i(x)|$$

$$= \text{In}\ b_i(s,t.x)\ b_i(t,x)$$

Additivity of $b_i(t,x)$ and the fact that $b_i(0,x)$ =0 implies

$$\frac{d}{dt}b_i(t,x) = \lim_{s\to 0}\frac{bi(t+s,x-bi(t,x)}{S-}$$

$$= \lim_{s\to 0}\frac{bi(s,x.t)}{S} = \frac{d}{ds}|_{s=0}\, b_i(s,x.t)$$

=

$$=(A_i(x.t)u_i(x.s),u_i(x.s),$$

The Fundamental Theorem of Calculus implies

$$B_i(t,x)= \int_0^t (A_i(x.s)\, u_i(x.s),\, u_i(x.s))ds\backslash$$

and the second assertion follows.

**Proposition 2:** *For the solutions of (2) we have:* if $x_i>0$, then $\lambda i(x) \le 0$. If $\lambda i(x)<0$, then $\lim_{t\to\infty} (x.t)_i=0$ If $limsup_{t\to\infty}$ $(x.t)_i|>0$, then $\lambda i(x)=0$.

**Proof:** First, assume that x satisfies $x_i>0$. Since the semi flow is dissipative, there exists $\gamma >0$ such that $|x.t| \le \gamma$ for all $t \ge 0$. Proposition 1 and the definition of the skew product flow imply

$$\lambda i(x) = \operatorname{limsup}_{t\to\infty}\frac{1}{t}\,\mathrm{In}\,|B_i(t,x)x_i| = \operatorname{limsup}_{t\to\infty}\frac{1}{t}\,\mathrm{In}\,|\,(x.t)_i\,|$$

$$\le \operatorname{limsup}_{t\to\infty}\frac{In\,\gamma}{t} = 0$$

To prove the second assertion, assume that $\lambda i(x) < 0. If\ x_i=0$, then the invariance of the faces of C imply that $(x.t)_i=0$ for all $t\ge 0$. Alternatively, if $x_i> 0$, then Proposition I and the skew product flow imply

$$\operatorname{limsup}_{t\to\infty}\frac{1}{t}In|(x.t)_i| = \lambda i(x) < 0.$$

The final assertion follows from the first two assertions. 5

**3.2. Invariant measures.** We review some definitions from ergodic theory. Given a Borel probability measure $\mu$ on ***C***, the *support* of $\mu$, denoted supp( $\mu$),is the smallest closed set whose complement has measure zero. A Borel probability measure $\mu$ is called *invariant* for (2) provided that Integral $h(x)d\mu(x) =$ Integral $h(x.t)d\,\mu(x)$for all

$t \geq 0$ and for all bounded continuous functions $h$:C→R. An invariant measure $\mu$ is called *ergodic* provided that $\mu(B)$=0 or 1 for any invariant Borel set *B.*

For an invariant measure $\mu$ define the invasion rate species $i$ with respect to $\mu$ as

$$\lambda_i(\mu) = \int_c \langle \mathrm{A}_i(x)\mathrm{u}_i(x), \mathrm{u}_i(x)\rangle d\mu(x)$$

**Proposition 3:** Let $\mu$ be an *invariant measure.* $\lambda_i(\mu)$ *Satisfies the following properties:* $\lambda_i(\mathrm{x}) = \lim_{t\to\infty} \frac{1}{t}\int_0^t \langle \mathrm{A}_i(x.s)\mathrm{u}_i(x.s), \mathrm{u}_i(x.s)\rangle, ds$ exists for $\mu$ - *almost every* $x$. *Moreover, if* $\mu$ *is ergodic,* $\lambda_i(\mathrm{x}) = \lambda_i(\mu)$ for $\mu$-*almost every* $x$.*If* $\mu$ *is ergodic, then there exists* $I \subset \{1, \dots, k\}$such that $\mu\,(\Pi_{i\in I}\boldsymbol{C}_i^+)$=1 and $\lambda_i(\mu) = 0\ for\ all\ \ i \in I$.

**Proof:** Let $\mu$ be an invariant measure. Define $\mathrm{h}_i(x) = \langle \mathrm{A}_i(x)\mathrm{u}_i(x), \mathrm{u}_i(x)\rangle$. The Birkhoff's' ergodic theorem and Proposition 1 imply

$$\lambda_i(x) = \lim_{t\to\infty} \frac{1}{t}\int_0^t \mathrm{h}_i(x.s)\mathrm{ds}$$

Exists for $\mu$- almost every $x$. Moreover, $\lambda_i(x) = \lambda_i(\mu)$ $\mu$- almost surely if $\mu$ is ergodic.

Assume $\mu$ is ergodic. By ergodicity and invariance of the faces of **C**, there exists $I \subset \{1, \dots, k\}$ such that $\mu\,(\Pi_{i\in I}\boldsymbol{C}_i^+)$=1. Let K $\subset \Pi_{i\in I}\boldsymbol{C}_i^+$ be a compact set such that $\mu(K) > 0$. The Poincare recurrence theorem and the Birkhoff ergodic theorem imply that there is $x \in \Pi_{i\in I}\boldsymbol{C}_i^+$such that $x.\mathrm{t}_n \in K$ for some $\mathrm{t}_n \uparrow \infty$ and $\lambda_i(x) = \lambda_i(\mu)$ for all $i$. The third assertion of Proposition 2 implies $\lambda_i(\mu) = 0$ for all $i \in I$.

**3.3. Robustly unsaturated sets**. Let K $\subset \partial C$ be a compact isolated invariant set for the flow $x.t$ restricted to $\partial C$.*K* is unsaturated if Ws (K) $\subset \partial C$ and K is isolated for $x.t$. If *K* is not unsaturated, then *K* is saturated. *K* is robustly unsaturated for (2) if there exists $\delta > 0$ such that the continuation of K for any $\delta -$ perturbation of (2) is unsaturated.**Theorem 1:** Let *K* be a compact isolated invariant set for $x.t$ restricted to$\partial C$. If one of the following equivalent conditions hold For all invariant measures $\mu$ supported by *K* (4) $\max_{1\leq i\leq k} \lambda_i(\mu) = 0$

There exist $\mathrm{p}_i > 0$ such that (5) $\sum_{1\leq i\leq k} \mathrm{p}_i\,\lambda_i(\mu) = 0$

For all ergodic probability measures supported by *K*.

then $K$ is robustly unsaturated for (2). Alternatively, if $x \to A(x)$ is twice continuously differentiable and $K$ is robustly unsaturated, then (4) holds for all ergodic measures $\mu$ supported by $K$.

**Remark 1**: For some applications (e.g., the disease model considered in section 6.1) it useful to relax the irreducibility assumption A2. For instance, if there exists an open neighborhood U of K such that for each $i$, $A_i(x)$ has a fixed off-diagonal sign pattern for all $x \in U$,then $A_i(x)$ can be decomposed into a finite number, say $m_i$, of irreducible components. For each of these irreducible components, one can define $\lambda_i^j(\mu) = \int_C \langle A_i^j(x) u_i^j(x), u_i^j(x) \rangle \, d\mu(x)$where $A_i^j(x)$ is the submatrix of $A_i(x)$ corresponding to the $j$-th irreducible component of $A_i(x)$ $and$ $u_i^j(x)$ is the continuous invariant sub bundle for the irreducible linear cocycle determined by $A_i^j(x)$. If we define $\lambda_i(\mu) = \max_{1 \le j \le m_i} \lambda_i^j(\mu)$, then all of the assertions of Theorem 1 still hold.

**Proof:** To prove the first assertion of the theorem, we need the following lemma. We call an invariant measure $\mu$ for (2) saturated if $\lambda_i(\mu) \le 0$ for all $i$

**Lemma 1**: Let $\delta_n$ be a non –negative sequence that converges to zero as n$\to \infty$. If $\mu_n$ are saturated invariant measures for $\delta_n -$ perturbations of (2), then the weak$^*$ limit points of $\{\mu_n\}_{n=1}^{\infty}$ is a non-empty set consisting of saturated invariant measures for (2).

**Proof:** Let $\delta_n$ be a non-negative sequence that coverages to zero. Let $\mu_n$ be saturated invariant measures for $\delta_n$ perturbations, $\dot{x} = A^n(x)x$, of $\dot{x} = A(x)x$. By weak$^*$ compactness of Borel probability measures supported on $N_1(\Gamma(A))$, there exist weak$^*$ limit points of $\mu_n$ . Let $\mu$ be such a weak$^*$ limit point. Since $\mu_n$ are supported by $N_1(\Gamma(A))$, for all $n$ sufficiently large, $\mu$ is supported by $N_1(\Gamma(A))$.To verify that $\mu$ is an invariant measure for (2),let $h: C \to R$ be a bounded continuous function. Let $t > 0$ and $\in > 0$ be given. Choose $n$ sufficiently large so that $|\int_c h(x)\, d\mu(x) - \int_c h(x) d\mu_n(x)| \le \in$, $|\int_c h(x.t)\, d\mu(x) - \int_c h(x.t) d\mu_n(x)| \le \in$, and $|h(x.t) - h(x^n.t)| \le \in$ for all $x \in N_1(\Gamma(A))$.Then $|\int_c h(x.t) - h(x) d\mu(x)|$ is

$\leq \quad |\int_c h(x.t)\,d\mu\,(x) - \int_c h(x.t)\,d\mu_n\,(x)| + |\int_c h(x.t)\,d\mu_n\,(x) - \int_c h(x)\,d\mu\,(x)|$

$\leq \quad \in + \int_c |h(x.t) - h(x^n.t)|\,d\mu_n\,(x)| + |\int_c h(x^n.t)|\,d\mu_n\,(x) - \int_c h(x)d\mu\,(x)|$

$\leq \quad 2 \in + |\int_c h(x^n.t) - h(x)\,d\mu_n\,(x)| + |\int_c h(x)d\mu_n(x) - \int_c h(x)d\mu\,(x)|$

$\leq \quad 3 \in$

Where the last line follows from the invariance of $\mu_n$ for $x^n.t$. Since $\in > 0$ and t > 0 are arbitrary, we have $\int_c h(x.t)d\mu\,(x) = \int_c h(x)d\mu\,(x)$ for all t >0 and all bounded continuous functions $h: C \to R$. It follows that $\mu$ is an invariant measure for (2). To see that $\mu$ is saturated, define $h_i^n(x) = \langle A_i^n(x)u_i^n(x), u_i^n(x)\rangle$ where $u_i^n(x)$ spans the invariant one-dimensional bundle given by $[26, Prop.\,3.2]$ for $\frac{dx}{dt} = A^n(x)x$. Since $A_i^n(x) \to A_i(x)$ uniformly for $x \in N_1((A))$ as $n \to \infty$, $[26, Prop.\,3.2]$ implies that $u_i^n(x) \to u_i(x)$ converges uniformly for $x \in N_1(\Lambda(A))$. Given $\in > 0$, choose $n$ sufficiently large so that $|h_i^n(x) - h_i(x)| \leq \in$ for all $x \in N_1(\Lambda(A))$ and $|\int_c h_i^n(x)\,d\mu\,(x) - \int_c h_i^n(x)\,d\mu_n\,(x)| \leq \in$.

Then

$\lambda_i(\mu) \leq \int_c |h_i(x) - h_i^n(x)|d\mu\,(x) + |\int_c h_i^n(x)\,d\mu\,(x) - \int_c h_i^n(x)\,d\mu_n\,(x)|$

$+ \int_c h_i^n(x)\,d\mu_n(x)$

$\leq \quad 2\in$

Since $\in > 0$ is arbitrary, $\lambda_i(\mu) \leq 0$ for all $i$ and $\mu$ is saturated.

Using this lemma, we prove that if K is saturated, then there exists a saturated invariant measure $\mu$ supported by K. Assume K is saturated. Work of Hofbaue and So $[7, Thm.\,2.1]$ implies that either there exists $y \in W^s(K) \cap C^+$ or K is not isolated for the unrestricted flow $x.t$ If there exists $y \in W^s(K) \cap C^+$, then for all $t \leq 0$ define $v_t = \frac{1}{t}\int_0^t \delta_{y.s}\,ds$ where $\delta_{y.s}$ denotes a Dirac measure based at the point $y.s$. Dissipativeness of (2) and weak* compactness of the Borel probability measures supported by (A) imply there exist $t_k \to \infty$ such that the sequence $v_{tk}$ converges in the weak* topology to a Borel probability measure $\mu$ with support in K.A standard argument

implies that $\mu$ is $x.t$ invariant. Define $h_i: C \to R$ by $h_i(x) = (A_i(x)u_i(x), u_i(x))$.Proposition 2 and weak* convergence imply that

$$0 \geq (\lambda i(y) = \limsup_{t\to\infty} \frac{I}{t}\int_0^t h_i(y.s)ds$$

$$= \limsup_{t\to\infty} \int_c h_i \, dv_t$$

$$\geq \limsup_{t\to\infty} \int_p h_i \, d\, v_{tk} = \lambda i(\mu)$$

For all $i$. Alternatively suppose that K is not isolated for the semi flow Then there exists a sequence of positive $\omega - limit$ sets that accumulate on K. Let $\mu_n$ be a sequence of ergodic probability measures supported by these $\omega -$ limit sets. Propositions 1 and 3 imply that $\mu_n$ are saturated for all $n$ Applying Lemma I with $\delta_n$=0 for all $n$ implies that there exists a saturated invariant measure $\mu$ suppoorted by *K*. Hence we have shown that if *K* is saturated then there exists a saturated invariant probability measure $\mu$ supported by *K* Equivalently, (4) holding for all invariant probability measures with support in *K* implies that *K* is unsaturated.

Next we show that if *K* not robustly saturated, then there exists a saturated invariant measure supported by *K* Indeed suppose *K* is not robustly saturated, Then there exists a non-negative sequence $\delta_n$ converging to zero and a sequence of saturated measures $\mu_n$ for $\delta_n$-perturbations of (2)with support in the continuation of *K*. Let $\mu$ be a weak* limit point of $\{\mu_n\}\infty_{n=1}$ Lemma 1 implies that $\mu$ is saturated. Moreover, since the continuation of *K* converges to *K* as $\delta_n \to 0, \mu$ is a saturated invariant measure for (2) supported by *K* Hence we have shown that if *K* is not robustly saturated then there exists a saturated invariant measure supported by *K* Equivalently if (4) holds for all invariant measures with support in *K* then *K* is robustly unsaturated.

To see the equivalence of the conditions given by (4) and (5) let $\Delta = \{p \in R_+^k : \sum_i P_i = 1\}$ and notice that

$$\text{Min max } \lambda i\ (\mu = \min_{\mu} \max_{p\epsilon\Delta} \sum_i Pi\lambda i(\mu)$$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K* The Minimax theorem (see, e. g.,[27] implies that

(6) $\min \max \lambda i(\mu) = \min_{\mu} \max_{p \epsilon \Delta} \sum_i Pi\lambda i(\mu)$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K*. Since min$\mu$ $\sum_i P_i \lambda_i(\mu)$ is attained at an ergodic probability measure with support in *K*, the equivalence of the conditions given by (4) and (5) is established.

To prove the final assertion of the Theorem assume $x \to A(x)$ is twice continuously differentiable and there exists a saturated ergodic measure $\mu$ supported by *K* Proposition 3 implies that there exists *I*C {1,…,k} such that (□$i \in iC_i^+$)=1and $\lambda i(\mu) = 0$ For all $i \in I. SinceKc\,.\partial$□,$\{1, \ldots . k\}\backslash I$ is non-empty. We will show that (2) is not robustly permanent by proving that for all $\delta > 0$ there exists a $\delta$ −perturbation of (2) for which the continuation of *K* is saturated Let $\delta > 0$ be given Choose V $c$ N$\delta$ (*Γ*(A)) to be a compact neighborhood of Γ(A) such that V.t $c$ int *V* for all t > 0.Let $\eta > 0$ be such that *W*=:*N* $\eta$(Γ(A))$c$ *V* Let p:c→ $[0,1] be$ a$C^{\infty}$function such that $p(x) = 1 \quad for\ all\ x \in \Gamma(A) and\ p(x) = 0 for\ all\ x \in C\backslash W\ Define\ \tilde{A} = (\tilde{A}_{1,}\ldots.,\tilde{A}_k)$ by

$$\tilde{A}_i(x) = \{ \frac{A(x)}{A(x) - \frac{\delta}{2Idp\,(x)}} \quad \frac{if\ i \epsilon I}{if\ i \notin I}$$

Where Id denotes the identity matrix of appropriate dimension Let x. t denote the semi flow of $x = \tilde{A}(x)x$.Since$x.t = x.t$ $whenever\ x.[0,t] \in C\backslash W, it\ follows\ that\ \ \tilde{V}.t \subset intV\ for\ all\ t >$ $0. Hence, \Gamma(\tilde{A})c\ V\ c\ N\delta\ (\ (A))). We\ \ also\ have\ ||A(x) - \tilde{A}(x)|| \leq \frac{\delta}{2}$Therefore $x = \tilde{A}(x) is\ a\ \delta -$ perturbation of (2) By construcation,$x.t = x.t\ for\ all\ x \in Kand\ t \geq 0. Hence, \mu$ is ergodic for the semiflow of $\dot{x} = \tilde{A}(x)x\ and\ \lambda_i(\mu) \leq -\frac{\delta}{2}$ for this semiflow and $i \in \{1, \ldots, k\}\backslash I$.Let $\mathcal{L}$ and O be the Lyapunov exponents and Oseledec regular points supported by $\mu$ for $x = \tilde{A}(x)x(see, e.g.,$

[28] for definitions).At each point $x \in O$, the splitting of $R^n$ determines three subspaces: the stable subspace $E^s(x)$, the center subspace $E^c(x)$, and the unstable subspace $E^u(x)$ Proposition 3 and our choice of Ã imply that $E^s(x) \cap C^+ \neq \emptyset$. The Pesin stable manifold theorem [28, Corollaries 3.17 and 3.18] implies that tangent to $E^s(x)$ $E^c(x)$ and $E^u(x)$ are locally $x.t-$invariant families of $C^1$ discs $W^s_x, W^c_x$ and $W^u_x$ corresponding to the stable ,center and unstable manifolds The family of stable manifolds $W^s_x$ is contained in $W^s$ (Ã,supp $(\mu)$). $Since\ E^s(x) \cap C^+ \neq \emptyset$, $W^s_x \cap C^+ \neq \emptyset$ for some $x \in K$. Consequently $K$ is saturated for $x = \tilde{A}(x)x$.

The difference between the sufficient and necessary condition is that for the sufficient condition (4) has to hold for all invariant measures supported by *K*, while for the necessary condition (4) has to hold only for ergodic measures supported by *K*. since the invariant measures lie in the convex hull of the ergodic measures the sufficient condition can be more restrictive then the necessary condition.

For three Classes of invariant sets the necessary and sufficient conditions coincide Recall a compact invariant set *K* for (2) is uniquely ergodic if K only supports one invariant measure Recall a compact invariant set *K* is Axiom A if the flow of (2) restricted to *K* is transitive and hyperbolic (see, e.g.,[29] for definitions.)

**Corollary 1:** Assume $x \to A(x)$ is twice continuously differentiable It K$\subset \partial$□ is a compact invariant set and either is uniquely ergodic or Axiom A, then *K* is robustly unsaturated if and only if (4) holds for all ergodic, measures $\mu$ with support in *K*

Proof. If K is uniquely ergodic then the assertion follows immediately For an Axiom A invariant set Sigmund [30, thm.1] has proven that invariant measures supported by periodic orbits of *K* are dense in the set of invariant probability measures supported by *K* In particular any invariant measure can be approximated by an ergodic measure and consequently the result follows.

Another special case where the necessary and sufficient conditions coincide is when *K* supports all populations except one for discrete time models this case was considered by [22].

**Corollary2:** Assume $x \to A(x)$is twice continuously differentiable Let K$c$ $II_i \geq 2C_i^+$ be a compact invariant set. Then K is robustly unsaturated if and only if $\lambda_1(\mu) > 0$ for all invariant measures supported by *K*

## 4. MORES DECOMPOSITIONS AND ROBUST PERMANENCE

To state the sufficient condition for robust permanence, we use a characterization of permanence due to Garay [8] and Hofbauer and So [7] that involves Morse decompositions of the boundary flow Conley[31] defined a collection of sets$\{M_{1,\ldots,}M_m\}$ to be a Morse decomposition for a compact invariant set $K$ if $M_1\ldots,M_m$ are pairwise disjoint, compact isolated invariant sent for the flow of (2) restricted to K. For each $x \in K$ there are integers r=r$(x) \leq s = s(x)$ such that $a(x) \subseteq M_r$and $\omega(x) \subseteq M_{s.}$ If $r(x) = s(x), then\ x \in M_{r.}$ Garay, Hofbauer and so [7,8] proved the following characterization of permamnce

**Theorem 2** (Garay, Hofbauer-so) if $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$?then (2) is permanent if and only if each of the components $M_1$ are unsaturated.

Theorem 1 and 2 imply the following result: As stated in the above all these results are true for most of problems, interactions, transactions in cosmology, econometrics, finance and other subjects, which we shall illustrate as we go on constructing models for the same.

**Theorem 3:** If $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$?and (4) holds for each of the components of the More decomposition, then (2) is robustly permanent Conversely if $x \to A(x) is\ twice\ continuously\ differentiable$

*and (4) is violated by an ergodic measure supported by one of the components of the Morse decomposition, then (2) is not robustly permanent.* Theorem 3 in conjunction with Corollaries 1 and 2 yield a characterization of robust permanence for a class of structured models.

**Corollary 3:** *Let $\{M_1,\ldots..,M_m\}$ os a Morse decomposition for* $\partial$C. *Assume x* $\mapsto$ *A(x) is twise continuously differentiable and for each Morse component $M_i$ one of the following assertions hold $M_i$ is Axiom A, $M_i$ is uniquely ergodic, or There exists j* $\in$ { 1,….,k} *such that $M_i$ ----- { $x \in C : x_l \gg 0$ for $l \neq j$} Then (2) is robustly permanent if and only if (4) holds for all ergodic measures supported by $U_iM_i$.*

Our result also provided a structured analogue for characterizing totally permanent systems[32]. For an ergodic probability measure$\mu$, define sup ($\mu$) to be the subset $I$ -----{1,…,k} such that $\mu(---_i \in_I C_+) = 1$

**Corollary 4:** *The following statements are equivalent: For all ergodic probability measures μ with support in ∂C, $\lambda_i(\mu)>0$ for all $I \in \{1,\ldots,k\} \setminus supp(\mu)$ (2) and all of its subsystems are robustly permanent.*

*Proof. Suppose the first statement holds. Then unsaturated condition given by (5) with p= (1,…,1) and Theorem 3 implies that (2) and all its subsystems are robustly permanent.. The other direction follows form the necessary condition for robust permanence in Theorem 3 and the second assertion of Proposition 3.*

*To illustrate the broad applicability of theorem 3, we develop applications to spatially structured models in section 5 and to a disease model and a gene network model in section 6.*

## 5. APPLICATIONS: PATCH MODELS

A fundamental application of our result are to spatially structured models, with *k* species dispersing between *m* patches {33-35}:

(7) $$\frac{dx_i^j}{dt} = x_i^j f_i^j(x^j) + \sum_l d_i^{jl} x_i^l - e_i^j x_i^j \qquad i = 1,..,k \qquad j=1,\ldots,m$$

Where $x_i^j$ denotes the density of species I in patch *j*, $x^j = (x^j_1,.....,x^j_k)$ is the vector of species densities in patch *j*, $f^j_i$ is the per-capita growth rate of species *i* in patch *j*, $d_i^{jl} \geq 0$ is the dispersal rate for species *i* from patch *l* into patch *j*, and $e_i^j$ is the emigration rate of species *i* out of patch *j*. Hence $e_i^j \geq \sum_l d_i^{lj}$. We assume that the matrices $(d_i^{lj})_{j,l}$ are irreducible for each species *i* . it is then easy to write (7) in the form (2), with $x_i = (x^l_i, \ldots, x_i^m)$ 'where' denotes transpose.

For a single species (7) generates a monotone flow. Under mild assumption, e.g., each $f^j_i$ *is* decreasing and negative for large densities, there is a globally stable equilibrium, see, e.g.,[36] and [37, sect. 5.4]. This equilibrium will be positive if the invasion rate at the origin λ (0), which is given by the leading eigenvalue (stability modulus) of the matrix $(d^{jl} + f^j(0) - e^j)$, is positive.

5.1 Two types of debitable and creditable systems:. For *two competing species*, e.g., $f^j_i$ is decreasing with respect to both species and is negative at large densities, the dynamics are still monotone, and hence almost all orbits converge

to an equilibrium. Robust permanence requires two conditions. First, $\lambda_i$ (0) >0 for both species in which case there are two single species equilibria $E_1$ and $E_2$. Second, the invasion rates $\lambda_2$ ($E_1$ ) and $\lambda_2$ ($E_2$) are positive. Whether there is a unique positive globally stable equilibrium for the robustly permanent system depends in a delicate way on the system parameters, see [38]. For slowly dispersing populations (i.e. $0< e^j_i \ll 1$ for all *i, j* ), the robust permanence condition is particularly

Straight forward to verify. Let $E^j_i$ be the largest solution to $x^j_i \quad f_i{}^j$ ($x^j_i$) = 0 i.e. the equilibrium attained by species *I* in patch *j* when there is no dispersal and no competitors. For slowly dispersing populations, $E_1$ is close to the product ($E_1^1$ ,…,$E_1^m$), and the matrix $A^2$ ($E_1$) is close to the diagonal matrix diag($f^i_2(E_1{}^j)$)$^m_{j=1.}$ Hence the invasion rate $\lambda_2$ ($E_1$) is close to $\max_j f_2^j (E^j_1)>0$. So we obtain: The system with two slowly dispersing competing species is robustly permanent if and only if system refers to system of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

(8) $\max_j f_2^j (E^j_1)>0$ and $\max_j f_2^j (E^j_2)>0$

Intuitively, robust permanence requires for each species there is at least one patch where it can persist.

For *systems, particle interactions, financial transactions, cosmological interactions and reactions, as we shall show model by model the* situation are even easier. Since (under mild assumption on the prey dynamics) the Morse decomposition of $\partial$C is given by two equilibria {0,$E_1$}, the spatial predator-prey system is robustly permanent, if the predator can invade at the prey equilibrium: $\lambda_2$ ($E_1$) > 0. However, the global dynamic are likely to be more complicated. For instance, if in each patch there is a globally stable limit circle (e.g.,[39-41]) and dispersal is sufficiently slow (i.e. $0< e^j{}_i \ll 1$for all *i, j*), our result about robust permanence and the theory of normally hyperbolic manifolds [42] imply that there is a positive *m* dimensional torus, which attracts almost all orbits in C, whenever $\max_j f_2^j (E^j_1)>0$

**5.2 Rock-Paper-Scissors Dynamics.** The Lotka-Volterra model of rock-paper-scissor dynamics is a simple model that is used as property for understanding intransitive ecological outcomes [43,44]. Here, a simple spatial version of this dynamic is given by

$\frac{dx_1^j}{dt} = x_1^j(1 - x_1^j - \beta^j\ x_2^j\ - \alpha^j\ x_3^j\ ) + \sum d_1^{jk}\ \ x_1^k\ - e_1^j\ x_1^j$

$\frac{dx_2^j}{dt} = x_2^j(1 - \alpha^j\ x_1^j\ - x_2^j - \beta^j\ \ x_3^j\ ) + \sum d_2^{jk}\ \ x_2^k\ - e_2^j\ x_2^j$

$\frac{dx_3^j}{dt} = x_i^j(1 - \beta^j\ x_1^j\ - \alpha^j x_2^j - \beta^j\ \ x_3^j\ ) + \sum d_3^{jk}\ \ x_3^k\ - e_3^j\ x_3^j$

Where $a^j \in (0,1)$ and $\beta^j > {}^1$ for all $j$. A more general version of this dynamic is presented in [44]. Under the assumption that there is no cost to dispersal (i.e. $e^j_i = \sum_l d^{l,j}_i$ for all $i, j$) , the maximal invariant set in $\partial C$ consist of the origin 0 and a heteroclinic cycle connecting positive single species equilibria$E_1$, $E_2$ and $E_3$. For slowly dispersing population ($0 < e^j_i \ll 1$ for all $i$, j), $E_1$ is close to $(x_1, x_2, x_3) = (1,\ldots,1,0,\ldots,0,0,\ldots,0)$ and the invasion rates $\lambda_2\ (E_1)$and $\lambda_3\ (E_1)$are close to $\max_j\ 1 - a^j$ and $\max_j\ 1 - \beta^j$. For the other equilibria $E_i$, the debitable rates of the missing species are also given by $\max_j\ 1 - a^j$ and $\max_j\ 1 - \beta^j$. Consider the Morse decomposition of $\partial C$ given by 0 and the heteroclinic cycle. Since $\lambda i\ (0) = 1 > 0$ for $i = 1, 2, 3, 0$ is robustly insaturated. An algebraic computation reveals that (4) holds for all invariant supported by the heteroclinic cycle (i.e. all convex combination of the Dirac measures supported by the equilibria) if and only if

$$\max_j(1 - a^j) > -\max_j\ (1 - \beta^j)$$

Equivalently,

$$2 > \min_j\ a^j + \min_j\ \beta^j$$

In particular, even if the heteroclinic cycle is attracting for each patch when the system is uncoupled (i.e. $2 < a^j + \beta^j$ for all $j$ ), it can be repelling for the weakly coupled system.

5.3. Three species **Lotka Volterra in spatially homogenous environments.** Consider Lotka-Volterra dynamics in a spatially homogenous environment: $(x + a)^n =$

(9) $$\frac{dx_i^j}{dt} = x_i^j\left(r_i - \sum_s a_{is}x_s^j\right) + \sum_l d_i^{jl}x_i^l - e_i^j x_i^j, \qquad i = 1,\dots,k;\ j = 1,\dots,m$$

We assume that dissipative can be shown by a linear Liapunov function: There are $c_i > 0$ such that $\sum_i c_i a_{is}x_i x_s < 0$ holds for all $x \in R_+^k \setminus \{0\}$. Then the weighted sum of densities across all patches $\sum_{i,j} c_i x_i^j$ is decreasing for large densitities and hence (9) generates a dissipative semiflow with global attractor $r$. We further assume that for each patch *j* and each species *i* the immigration rate equals the emigration rate , i. e,

(10) $$e_i^j = \sum_l d_i^{jl}$$

This guarantees that the set $H = \{x \in C: x_i^j = x_i^l \forall i,j,l\}$ of spatially homogenous states is forward invariant under (9).

Now consider the Lotka-Volterra dynamics without spatial structure

(11) $$\frac{dx_i}{dt} = x_i\left(r_i - \sum_s a_{is}x_s\right), \qquad i = 1,\dots,k$$

(11) $$\frac{dx_i}{dt} = x_i(r_i - a_{is}x_s), \quad i = 1,\dots,k$$

It has been shown [12,ch.16.1,16.2] That for $k = 3$,(11) is robustly permanent (robust meaning here within the class of Lotka –Volterra systems) if and only if all equilibria on the boundary of $R_+^3$ are unsaturated and whenever there is a heteroclinic cycle connecting the one species equilibria then this cycle is repelling. Equivalently the following four conditions hold. There exists an interior equilibrium $\dot{x}$ (i.e $A\dot{x} = r \ \ with\ \dot{x} \gg 0$); Det (-A)>0; the 2 species subsystems are not bistable competition systems if there is a heteroclinic cycle between the one species equilibria, say $\mathbf{E_1} \to \mathbf{E_2} \to \mathbf{E_3} \to \mathbf{E_1}$ then the following inequality holds

12) $$\lambda_2(E_1)\lambda_3(E_2)\lambda_1(E_3) > |\lambda_3(E_1)\lambda_1(E_2)\lambda_2(E_3)|$$

Furthermore, the proof in [12, ch.16.1, 16.2] shows that the boundary flow of (11) has a simple Morse decomposition: if (iv) applies then there are two Morse sets, the origin 0 is a repeller, and the heteroclinic cycle is the dual attractor (within$\partial R^3_+$); if (iv) does not apply then the finitely many boundary equilibria form a Morse decomposition of the boundary flow.

**Species, population we mean types of and population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.**

**Theorem 4:** Under the above assumptions,(9) for k=3 is permanent if and only if *(i),(ii),(iii)* and *(iv)* hold.

Proof: Suppose (9) is robustly permanent. Since the nonspatial system (11) is the restriction of (9) to the invariant subspace **H**, (11) must be robustly permanent and hence*(i)-(iv)* hold

Conversely, suppose that (i)-(iv) hold. We show that the above Morse decomposition of the boundary flow of (11) (now in **H**) is also Morse decomposition for (9) restricted to the maximal invariant subset of $\partial C$ which is contained in the spatial two species subsystems of (9). For this we show that in each 2 species system the orbits of (9) converge to **H** (For the one species subsystems this is obvious.) If a two species subsystem of (11) has an internal equilibrium, say $\mathbf{E_{12}}$ then by assumption (ii),this is not a saddle but a sink, and hence $a_{11}a_{22} > a_{12}a_{21}$. This implies that the 1-2 submatrix of A is VL-stable [12,ch.15.3,],and [45] implies that $\mathbf{E_{12}}$ is the global attractor for the spatially structured system (9) restricted to the first two species.

If a two species subsystem of (11) has no internal equilibrium, then a one species equilibrium say $\mathbf{E_1}$ is the global attractor. Then Lemma 2 (see below) show the corresponding result for (9).A similar proof applies to the case where the origin **0** is the global attractor of a two species subsystem.

Finally observe that under the above assumptions (9)-(10), the invasion rates $\lambda_i$ at spatially homogenous boundary equilibria are the same for (9) and (11) Since all boundary equilibria are unsaturated for (11), they are for (9) as well. The same applies in case (iv) to the heteroclinic cycle. Note that the condition (12) on the eigenvalues is equivalent to the fact that all invariant measures supported on the heteroclinic cycle are unsaturated, see [16,Ex.4.5]

Species we mean types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

**Lemma 2:** Suppose, in a two species Lotka –Volterra system (11) (k=2) all interior orbits converge to the one species equilibrium $\mathbf{E_1}$ .Then all orbits of the spatial version (9) (k=2) converge to the spatially homogeneous one species equibrium $\mathbf{E_1} \in \mathbf{H}$.

**Proof:** If the system is of predator-prey or mutualistic type then Hasting's [45] result applies again. So let us assume that the local interaction is given by the competition system

(13) $\frac{dx_1}{dt} = r_1 x_1(1 - x_1 - ax_2)$ $\frac{dx_2}{dt} = r_2 x_2(1 - \beta x_1 - x_2)$

With 0<a <1< $beta$.In this case, Kazuo Kishimoto (letter to JH, Oct 1987) has given the following argument. There exists a family of forward invariant rectangles contracting to $\mathbf{E_1}$. This shows that all interior solutions of (9) (k=2) converge to $\mathbf{E_1}$.

Theorem 4 shows that a spatial network of identical patches is permanent, if the within patch dynamics is permanent, and three species Lotka-Volterra dynamics. It is not clear how to extend this result to more than three species. The crucial step in the three species case is that permanence precludes bistable two species subsystems. The spatial version of a bistable two species system allows plenty of stable equilibria outside **H,** see [33].However, permanent four species Lotka Volterra systems may have two species bistable subsystems [12,ch.16.4].So there may be invariant sets on the boundary outside **H** which need to be unsaturated to make (9) permanent.

## 6. MORE APPLICATIONS

6.1. **Disease, particle, product, marketing, financial, econometric dynamics with density-dependent demography**. GAO and Hethcote [46] introduced a model of disease dynamics with density-dependent demography. Here we consider a variation of their model in which a population of size P has a per-capita birth rate b(P) and per –capita death rate d. To ensure that the population persists in the absence of the disease, we assume that there exists K>0 such that b(K)=d. To describe the disease dynamics, let *S* be the number of individuals susceptible to the disease, *I* the Number of infected individuals, and R the number of individuals that have recovered from the disease. We assume that P=S+*I* +R. On the same lines we can think of sick and non performance accounts, and those which are conducting themselves well, and those which are satisfactory and find a parallel formula. The same thing

holds true in case of a black hole eating matter and black hole radiating radiation.. If the disease transmission is asymptotic, β is the contact rate between susceptible and infected individuals, γ is the rate at which individuals recover from the disease, and m the mortality rate due to disease, then the population dynamics is given by

$$\frac{dp}{dt} = b(P)P - dP - mI$$

$$\frac{dI}{dt} = \beta(P - I - R)\frac{I}{\epsilon + P} - (\gamma + d + m)I$$

$$\frac{dR}{dt} = \gamma I - dR$$

Where $\epsilon>0$ is a constant It is useful to introduce a change of coordinates in which $y=\frac{I}{P}$ and $z=\frac{R}{P}$:

$$\frac{dP}{dt} = (b(P)-d-my)P$$

$$\frac{dy}{dt} = \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(P)\ y$$

$$\frac{dz}{dt} = \gamma y - (b(P) - my)z$$

Setting $x_1$=P and $x_2$=(y,z) 'where' denotes transpose yields a structured model where $A_1(x)$=b-d-my and

$$A_2(x) = \begin{pmatrix} \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(p) & O \\ \gamma & my-b(p) \end{pmatrix}$$

Since {0,($K$,0,0)}is a Morse decomposition of the boundary dynamics Theorem 3 and Remark 1 imply that this model is robustly permanent if and only if b(0)>d and $\beta\frac{K}{\epsilon+K} > \gamma + m + d$. In particular, for $\epsilon > 0$ sufficiently small, one requires that the basic reproductive number $\frac{\beta}{r+m+d}$ is greater then one.

Remarkably this same criterion determines robust permanence of model with significantly more complicated boundary dynamics. For example, suppose the focal population is a predator species. If the prey has abundance $N$ and exhibits logistic dynamics f($N$)is the per-capita predator consumption rate of the prey and b(N)is the per-capita reproductive rate of the predator then the dynamics become

$$\frac{dN}{dt} = rN(1 - N/K - f(N)P$$

$$\frac{dP}{dt} = (b(N) - d - my)P$$

$$\frac{dy}{dt} = \left[ \beta(1 - y - z)\frac{P}{\epsilon+P} - r - m(1 - y) - b(N)\ y \right]$$

$$\frac{dz}{dt} = \gamma y - (b(N) - my)z$$

Where r>0 and *K*>0, and *f*(*N*),b(N)satisfy f(0)=b(0)Define $x_1$=*N*, $x_2$=*p*, $x_3$=*(y,z)*′ Let us assume that the disease-free community (i.e.y=z=0)is permanent (i.e.b(K)>d) and let A be the global attractor in the interior of the P-N plane Under these assumption a Morse decomposition of the boundary dynamics is given by $(x_1, x_2, x_3) = (0,0,0)$, $(x_1, x_2, x_3) = (K, 0,0)$and A. Since the equilibria (0,0,0) and (k,0,0) are unsaturated, it remains to characterize $\lambda_3(\mu)$for any invariant measure μ supported by A. Proposition 3 and Remark 1 imply

$$\lambda_2(\mu) = \int_c (bx_1)\text{-}d)d\mu(x) = 0$$

*Hence* $= \int_c bx_1)d\mu(x) = d$ *and*

$$\lambda_3(\mu) = \beta \int \frac{x_2}{\epsilon+x_2} d\mu(x) - r - m - d$$

**Theorem 3**: and Remark 1 implies that this structured model is robustly permanent if $\frac{\beta}{r+m+d} > 1 and\ \epsilon > 0$ is sufficiently small.

6.2. Gene networks. The repressilator is an oscillatory gene network based on three (or more generally an odd number of transcriptional repressors. Such a system has been genetically engineered in E-coli by Elowitz and Leibler [47] Mathematical model go back to [48.49] they all involve concentrations of proteins and mRNs. A modified model with auto-activation was suggested in [50]

$$(14)\ x_i = \beta(y_i - x_i)$$

$$(15)\ y_i = aF(x_{)i} \text{-} y_i)$$

With

$$(16)\quad F(x)_i = x_i g\ (x_i, x_i - 1) = \frac{x_i}{1+x_i+p x_{i-1+Kp\, x_i x_i - 1}}$$

Here $x_i$ and $y_i$ are normalized concentrations of proteins and mRN As belonging to gene $i$,p is the strength of repression and k a cooperativity parameter this is a structured system of type(2)

$$(17)\ \begin{pmatrix} x_i \\ y_i \end{pmatrix} dot = \begin{pmatrix} -\beta & \beta \\ ag(x_i,x_i) & -1 \end{pmatrix} \begin{pmatrix} x_i \\ y_i \end{pmatrix}$$

In [50,Thm 3] it was shown that for n=3,a>1 and p>1,system(14)has a heteroclinic cycle connecting the 3 single gene equilibria (similar to the rock –scissors-paper dynamics of section5.2)Moreover

- the system is permanent i.e. the boundary of $R^6$+is repelling, if $\lambda + \mu > 0$
- this heteroclinic cycle is asymptotically stable if $\lambda + \mu < 0$

where $\lambda > 0$ and $\mu < 0$ are the invasion rates of the two missing genes i.e. the leading eignvalues of the 2X2 matrix in (17)evaluated at the appropriate single gene equilibria (for explicit expressions see[50,eqs.(127,128)].)

Theorem 3 implies the permanence condition implies the stronger conclusion of robust permanence A similar result holds for n odd, but the heteroclinic cycle is then actually a heteroclinic network.

Note: For references please see the original paper.

One of the examples for which we apply the model are as follows: **Progress in Particle and Nuclear Physics Volume 88, May 2016, Pages 1–28 Chiral magnetic and vortical effects in high-energy nuclear collisions—A status report D.E. Kharzeeva, b, J. Liaoc, d, , , S.A. Voloshine, G. Wangf**

The interplay of quantum anomalies with magnetic field and vorticity results in a variety of novel non-dissipative transport phenomena in systems with chiral fermions, including the quark–gluon plasma. Among them is the Chiral Magnetic Effect (CME)—the generation of electric current along an external magnetic field induced by chirality imbalance. Because the chirality imbalance is related to the global topology of gauge fields, the CME current is topologically protected and hence non-dissipative even in the presence of strong interactions. As a result, the CME and related quantum phenomena affect the hydrodynamical and transport behavior of strongly coupled quark–gluon plasma, and can be studied in relativistic heavy ion collisions where strong magnetic fields are created by the colliding ions. Evidence for the CME and related phenomena has been reported by the STAR Collaboration at Relativistic Heavy Ion Collider at BNL, and by the ALICE Collaboration at the Large Hadron Collider at CERN. The goal of the present review is to provide an elementary introduction into the physics of anomalous chiral effects, to describe the current status of experimental studies in heavy ion physics, and to outline the future work, both in experiment and theory, needed to eliminate the existing uncertainties in the interpretation of the data. Keywords

Chiral magnetic effect; Chiral vortical effect; Chiral anomaly; Quark-gluon plasma; Heavy ion collisions Corresponding author at: Physics Department and Center for Exploration of Energy and Matter, Indiana University, 727 E Third Street, Bloomington, IN 47405, USA. Copyright © 2016 Elsevier B.V. All rights reserved.

## Sufficient Conditions For The Existence And Linear Stability Of Time Independent Age Distributions In Respect Of Particles, Propositions, Problems, Hegelian Structures, Deleuzean Treatises, Interactions And Transactions In Cosmos And Others

**Abstract:** Gopalaswamy studied sufficient conditions are obtained for the existence and linear stability of time independent age distributions in two species competition with age time lagged density dependent mortality and fertility functions.. Work was published in Bulletin Australian Mathematical Society Volume 25 (1982). We show how the model could be applied in case of finance, econometrics, psycho economics, interactions between particles and transactions in nature. Basically the competition is for resources, notwithstanding such experiments could be conducted in laboratory like a rat kept in front of a cat, such incidents could be represented mathematically, and the corresponding linear stability conditions ca be obtained in case of interactions in cosmos. Take for example the case of a cash credit account of a borrower. Many inward bills waiting for paid, or inland or foreign bills payment all are in line for the debit of the account, which also acts as a resource. In addition there are recurrent cases where in the credit is issued to the cash credit account in case of demand and inland bills. In case of economics, investment has inventories and it has to be utilized towards the end of consummation of output. And this is case of output gobbling part of investment which is inventories. Output has to be sold to the customer and the amount received is diverted partly to savings and consumption and part could be employed towards the consolidation of the system. So this is a case of investment itself acting as a both receiver and payer. It shall eat and also produce the finance. Savings as stated above could be deployed for the further diversification and proliferation and that is towards the end of putting the company on terra firma. Another example in cosmology is the case of black holes guzzling matter. There are lot of black holes and lot of matter. Result of such a guzzling the production of Hawking radiation. Horava gravity is another example. How the conformal field theory could be changed in the case of Horava gravity. Still another case is the case of Vonneumann algebra in the cae of Horava gravity. In this case like cash credit account, the Horava gravity and Von Neumann algebra act as both receivers and takers. Such incidents are legion in cosmology and even in philosophy. We have considered self as individual general ledger, and the debits and credits are raised to the same account thus anointing as both a receiver and payer. We give plethora of such cases and apply the model. Due to spatial constants, same module is used for all the modules and we rewrite to state that the model is applicable to all modules. This is due to spatial restrictions. What is to be done is that each and every parameter is to be changed with a subscript or a superscript and finally it is to be concatenated towards the end of obtention of prediction model or optimizing model, for Duality of motion model for further analysis. Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation

of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

**Introduction**

Let $P_1$ *(t)* and $P_2$ *(t)* denote the total particles, interactions, transactions, etc, which we dilate upon further investigation. Perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. **Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems** sizes (or biomasses) at time $t \geq 0$ of two interacting species living in a common habitat and competing for a common pool of resources. The competition for resources will be implicit in our model similar to that in the two species Lotka- Volterra competition system. Assuming constant sex ratios in the two species we can consider $P_1$ *(t)* and $P_2$ *(t)*to be the Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

population of f; immigration, emigration and internal dispersion in the habitat are assumed to play no significant role in the dynamics of the community.Popoulation we mean Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

We suppose that the two species contain respectively $p_1(a, t)\,da$ and $p_2(a, t)\,da$ individuals with ages between $a$ and $a + da\ (a \geq 0)$ at time $t$ so that we have

(1.10) $P_1(t) = \int_0^\infty \rho_1(a, t)\,da,\ P_2(t) = \int_0^\infty \rho_2(a, t)\,da,\ t \geq 0.$

The rates of change of the two species densities are by definition $D\rho_1(a, t)$ and $D\rho_2(a, t)$ where

(1.11) $D\rho_i(a, t) = \limsup_{h \to 0} \frac{\rho_i(a+h, t+h) - \rho_i(a,t)}{h},\ i = 1, 2.$

Assuming the existence of the limits in (1.2) we consider the following time lagged model system

(1.12) $$\begin{cases} D\rho_1(a,t) + f_1(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_1(a, t) = 0, & a>0,\ t>0 \\ D\rho_2(a,t) + f_2(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_2(a, t) = 0, & \text{(1.4)} \\ \rho_1(0, t) = \int_0^\infty b_1(a, P_1(t), P_2(t-\tau))\rho_2(a, t)da, & t>0, \\ \rho_2(0, t) = \int_0^\infty b_2(a, P_1(t), P_2(t-\tau))\rho_1(a, t)da, & \end{cases}$$

(1.5) $\rho_i(a, s) = \varphi_i(a, s),\ a.0,\ s \in [-\tau, 0],$ $i = 1, 2,$

(1.6) $P_i(s) = \int_0^\infty \varphi_i(a, s)\,da,\ s \in [-\tau, 0],$

Where $\tau$some is fixed nonnegative constant; $f_1$ and $f_2$ denote the age and density dependent mortality rates with time lags in interspecific interactions, $b_1$ and $b_2$ denote the age and density dependent fertility functions again having time lags in the interspecific interactions; $\varphi_1$ and $\varphi_2$ denote the initial age distributions needed for the formulation of the model. Note we mean the population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

The model system (1.1)-(1.6) will be meaningful only if $\varphi_1$, $\varphi_2$, $f_1$, $f_2$, $b_1$, $b2$ are nonnegative functions of the respective arguments. Also since the total initial populations have o be finite, $\varphi_1(\cdot, s)$ should belong to $L_1(R^+)$ for each $s\epsilon$ $[-\tau, 0]$. That $f_i$ and $b_i$ $(i = 1, 2)$ are independent of t explicitly and this assumption can be interpreted to represent the temporal uniformity of the environment.

One of the fundamental questions for (1.1)-(1.6) is the following; under what conditions on $f_i$, $b_i$ and $\tau$ does the system (1.1)-(1.4) have time independent nontrivial solutions, and when such solutions, and when such solutions exist, are they stable with respect to some suitable stability criterion? When such stable time invariant solutions exist for (1.1)-(1.4) we say that the two types has a stable age distribution. Existence of a stable age distribution is the analogue of the existence of a stable steady state for the corresponding age independent system. Types we mean types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

In this article we investigate the above question for the case of a competitive interaction with some additional constraints on the vital rates $f_i$ and $b_i$ $(i = 1, 2)$. Age dependent population systems without time delays have been considered by several authors (Gurt in and MacCamy [3, 4, 5], Gurt in and Levine [6], Haimovici [7] and Rotenberg [8]). The question of the relation between an age independent system and a corresponding age dependent system has been considered by Gurt in and MacCamy [4] and Gurt in and Levine [6] who have established an asymptotic (*as* $t\rightarrow\infty$ ) relation between such models by constructing a higher dimensional 1 umped parameter system (in terms of ordinary differential equations) to represent the age dependent distributed parameter system. Assuming the existence of stationary age distributions Haimovici [7] considers their stability in a system of two interacting populations explicitly taking into consideration the dynamical nature of the encironments resources and pollution.

In this article we add another realistic feature namely time lags in the age dependent models and obtain sufficient conditions for the existence of stationary age distributions in (1.1)-(1.4) and show that the time lags in interspecific

interactions have no effect on the linear stability of the age distributions although the decay rates of perturbations will depend on the magnitudes of the time delays. Such "harmless" natures of time lags in interspecific interactions have been noted by the author for the age independent models in [1, 2].

## 4. Existence and uniqueness of solutions

We first narrate our assumptions on the vital rates of $f_i$ and $b_i$ $(i = 1, 2)$ so that the system (1.1)-(1.6) will denote a competition system. Let $BC\ (0, \infty)$ denote the linear space of bounded continuous real functions on $(0, \infty)$ with the norm defined by

$$||f|| = \{max\ |\ f(t)\ |,\ t \in (0, \infty)\}.$$

Let $K_1$ be the cone of nonnegative functions of $BC\ [0, \infty]$. Let $C\ [-\tau, 0]$ denote the space of real continuous functions with a norm defined by $||g|| = \{\ max\ |g(s)\ |,\ s \in [-\tau, 0]\}$; let $K_2$ be the nonnegative cone of functions in $C\ [-\tau, 0]$ . We present our first set of assumptions on $f_i$, $b_i$ and $\varphi_i$ for $i = 1, 2$.

$A_1$. $f_1 \in C\ (R^+ \times K_1 \times K_2, R^+)$, $f_1 \geq d_1^*$ on $R^+ \times K_1 \times K_2$

$f_2 \in C\ (R^+ \times K_2 \times K_1, R^+)$, $f_2 \geq d_2^*$ on $R^+ \times K_2 \times K_1$ ,

( $d_1^*$ , $d_2^*$ are positive constants);

$f_1$ and $f_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ ; that is, *for* $i = 1, 2$,

$$|f_i\ (a, P_1^{'}\ , P_2^{'}) - f_i(a, P_1^{"}\ , P_2^{"})| \leq d_i\ [||P_1^{'} - P_1^{"}, P_2^{'} - P_2^{"}||]$$

($d_1$, $d_2$ are positive constants),

$$\partial f_1/\ \partial P_1, \partial f_1/\ \partial P_2 \in C(R^+ \times K_1 \times K_2, R^+)\ ,$$

$$\partial f_2/\ \partial P_1, \partial f_2/\ \partial P_2 \in C(R^+ \times K_2 \times K_1, R^+),$$

$A_2$. $b_1 \in C\ (R^+ K_1 \times\ K_2, R^+)$, $b_1 \leq \beta_1^* on\ R^+ \times K_1 \times K_2$ ,

$b_2 \in C\ (R^+ K_2 \times\ K_1, R^+)$, $b_2 \leq \beta_2^* on\ R^+ \times K_2 \times K_1$ ,

*($\beta_1^*$ , $\beta_2^*$ are positive constants);*

$b_1$ , $b_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ *; that is for* $i = 1, 2$,

$$|b_i\ (a, P_1^{'}, P_2^{'}) - b_i\ (a, P_1^{"}, P_2^{"})| \leq \beta_i\ [||P_1^{'} - P_1^{"}|| + ||P_2^{'} - P_2^{"}]$$

$B_1$ , $B_2$ Are positive constants),

$b_i(°,(P_1,P_2)\epsilon L^{\infty}(R^+,R^+)$ For fixed $P_1,P_2,i=1,2,$

$$\frac{\partial b_1}{\partial P_2},\frac{\partial b_1}{\partial P_2}\epsilon\ C\ R^{+}\text{x}\ k_{\ 1}\text{x}\ k_{\ 2},\text{R}^{-}),$$

$$\frac{\partial b_2}{\partial P_1},\frac{\partial b_2}{\partial P_2}\epsilon\ C\ R^{+}\text{x}\ k_{\ 2}\text{x}\ k_{\ 1},\text{R}^{-}),$$

$A_3$. $\varphi_i\ (\cdot\ ,\text{s})\in \text{L}^{\ 2}\ (\text{R}^{+},\text{R}^{+})\ \cap \text{C}(\text{R}^{+},\text{R}^{+})$ for each $s\in\ [-\tau,0]$,

$i,\ =1,\ 2,$

By a solution of $(1.1)-(1.6)$ on $[0,\text{T}]$, we mean

$\text{P}_i$: $\text{R}^{+}$ ×$[-\tau,\text{T}]\rightarrow \text{R}^{+}$ $(i,\ =1,\ 2,)$ With the following properties:

$S_1$. $\text{P}_i(\cdot\ ,t)\in \text{L}^{\ 1}\ (\text{R}^{+},\text{R}^{+})$ for each $t\ \in[-\tau,\text{T}]$;

$S_2$. $\text{P}_i(a,\cdot\ )\in \text{C}\ ([-\tau,\text{T}],\text{R}^{+})$ for each $a\ \in \text{R}^{+}$;

$S_3$. $\text{DP}_i$ Exists along the characteristics $t-a=$ constant

On $\text{R}^{+}$ ×$\{\text{R}^{+}\cap\ [0,\text{T}]\}$ And is continuous $(i,\ =1,\ 2,)$;

$S_4$. $\text{P}_i$ Satisfies the system (1.1)- (1.6) for

$(a,t,)\ \in \text{R}^{+}$ ×$[-\tau,\text{T}]$ .

Let us first convert the system (1.1)- (1.6) into an equivalent system of integral equations. For brevity we denote by P the pair ( $\text{P}_1,\text{P}_2$) If we let

$$\text{P}_i(a+x,t+x)=\ \ \text{P}_i(x),(i,\ =1,\ 2,)\ ,$$

Then (1.3) considered at $(1.3)$ considered at $(a+x,\ t+x)$ becomes

$$\frac{dp_{\ i}(\text{x})}{dx}+f_i(a+x,\ P(t+x))\ \ P_i(x)=0$$

Which has a solution of the form?

$$P_i(x)=P_i(0)\ \exp\left[-\int_0^{\infty} f_{\ i}\big(a+s,P(t+S)\big)ds\right]$$

And hence

$$(2.1)\ p_{\ 1}\ (a+x,t+x)=p_{\ 1}\ (a,t)\ exp\big[-\int_o^x\ f_{\ 1}\ \big(a+s,p(t+s)\big)ds\big]$$

For $a \geq t$ *we choose (a,t)=(a-t,0)* in (1.3) and *x=t* in (2.1) so that

(2.2) $p_1(a,t) = \varphi_1(a-t,0)\exp\left[-\int_o^t f_1\left(a-t+s, p_1(s)\right)ds\right]$

$$O \leq t \leq a, i = 1,2,$$

With

$$P_1(s)=(p_1(s),p_2(s-t)),$$

$$P_2(s)=(p_1(s-t),p_2(s)),$$

For $0 \leq a < t$ we choose *(a,t)=(0,t-a)in (1.3)and x=a in (2.1)* so that

*(2.3)* $p_i(a,t)= B_i(t-a)exp[-\int_o^a f_1(s, p_i(t-a+s))ds]$

$$t > a \geq, i = 1,2,$$

Where

$$(2.4)\ B_i(t) = p_i(0,t), i = 1,2, t > 0.$$

Define $M_i and\ L_i\ (i = 1,2)$ as follows:

$$M_i(a,t,p_i(t)) = exp[-\int_o^{t-a} f_1(s, p_i(a-a+s))ds]$$

$$(2.5)\ L_i\left(a,t,\ p_i(t)\right) = \exp[-\int_o^t f_1(\text{a+s},\ p_i(s))ds]$$

It will now follow from (1.1),(1.3),(2.2)-2.5)that

$$(2.6)\ p_i(t) = \int_o^t B_1(a)\, M_1(a,t,p_i(t)da + \int_o^\infty \varphi_i(a,0)\ L_i(a,t,p_i(t)da,$$

$$t > 0, i = 1,2,$$

$$(2.7)\ p_i(s) = \Phi_i(s) = \int_o^\infty \varphi_i(a,s)da, s \in [-\tau, 0], i = 1,2,$$

(2.8) $\beta_i(t) = \int_0^t \beta_i(a) b_i(t-a, P_i(t)) M_i.\left(a,t,P_i(t)\right)da$

$$+\int_0^\infty b_i\left(a+t, P_i(t)\right)\varphi_i(a,0)L_i\left(a,t,P_i(t)\right)da$$

$$i = 1, 2, t > 0$$

the equivalence of (2.6)-(2.8) with (1.1)-(1.6) is established by the following whose proof is identical to a similar result of Gurtin and MacCamy [4]; hence we omit the details of the proof of the following.

THEOREM 1 : *Suppose the system* (1.1)-(1.6) *satisfies the conditions* $A_1, A_2, A_3$ . *If* $(P_1, P_2)$ *is a solution of*

*(1.1)-(1.6) then the total population* $P_1, P_2$ *and the birth rates* $B_1, B_2$ *satisfy the integral equation (2.6)-(2.8).*

*conversely if* $P_1, P_2, B_1, B_2$ *are nonnegative continuous solution of (2.6)-(2.8) on [ 0, T] and if* $P_1, P_2$ *are*

*defined by (2.2)-(2.4) on* $R^+$ x [0, T] *then such* $P_1, P_2$ *provide a solution of (1.1)-(1.6) on* $R^+$ x [0, T] .

The following a *priori* estimates are useful to prove our existence theorem below. It will immediately follow from the bounds for $f_i$ and $b_i$ that

$$M_i(a,t,P) \leq \exp[-d_i^*(t-a)] \leq 1, if \quad t \geq a,$$

*(2.9)*

$$L_i(a,t,P) \leq \exp[-d_i^* t]$$

From (2.7)-(2.8) we derive that

$$\beta_i(t) \leq \int_o^t \beta_i^* \beta_i(a) \exp[-d_i^*(t-a)]\, da + \beta_i^* \Phi_i(0) \exp[-d_i^* t], \quad i = 1,2,$$

And hence by Gronwall's inequality,

(2.10) $$\beta_i(t) \leq \beta_i^* \Phi_i(0) \exp[\delta_i t], \ \delta_i = \beta_i^* - d_i^*$$

(2.5), (2.6) and (2.10) lead to

(2.11) $$|P_i(t) - 0_i(0)| \leq \beta(0) exp[\delta_i a - d_i^*(t-a)]\, da$$

$$+\int_0^\infty \varphi_i(a,0)|1 - exp(-d_i^* t)|da$$

$$\leq \beta_i^* \Phi_i(0) exp(\beta_i^* t)\left\{\frac{1-exp[-(d_i^*+\beta_i^*)t]}{d_i^*+\beta_i^*}\right\} + \beta(0) d_i^* t$$

$$\leq [\beta_i^* \beta(0) exp(\beta_i^* t) \beta(0) d_i^*] t \quad for \ t > 0.$$

We can now prove the following result on the existence and uniqueness of solution of the integral equations (2.)-(2.8)

**THEOREM 2:** *Assume that the system* (1.1)-(1.6) *satisfies the hypotheses*

$A_1, A_2, A_3$. *Then there exists a unique set of continuous function* $(P_1, P_2, B_1, B_2)$ such *that*

$$P_i : R^+ \to R^+, \qquad B_i : R^+ \to R^+ \ (i = 1,2),$$

*Which satisfy the integral equations* (2.6)-(2.8)

Proof. We will first show the existence of local solution on same (possibly small) interval [0, T] $\subset R^+$ such that

$$P_i \in \mathrm{C}([0,T], R^+), \quad B_i \in \mathrm{C}([0,T], R^+), i = 1,2,$$

Satisfying (2.6)-(2.8) and then show the $P_i$ and $B_i$ can be continued as solution for any finite interval in $R^+$ for some fixed

Positive number $T$ we let

$$C^+[-\tau, T] = \{f \in \mathrm{C}([-\tau, T], R) | f \geq 0\}$$

Let $||.||_T$ denote the supremum norm in $\mathrm{C}^+[-\tau, T]$. Then for any

$(x, y) \in \mathrm{C}^+[-\tau, T] \times \mathrm{C}^+[-\tau, T]$ we define

$$||(x,y)|| = ||x||_T + ||y||_T,$$

$$\mathrm{C}_T = \mathrm{C}^+[-\tau, T], \quad \mathrm{C}_T^2 = \mathrm{C}^+[-\tau, T] \times \mathrm{C}^+[-\tau, T].$$

Members of $\mathrm{C}_T$ will be denoted by $Z_t$ with the meaning that

$Z(t+s) = z_t(s)$ *for* $s \in [-\tau, T]$ *and* $t \in [0, T]$. Now for any fixed

$P_t = (P_{1t}, P_{2t}) \in \mathrm{C}_T^2$, (2.8) is a linear system of uncoupled Volterra integral equations in B = $(B_1, B_2)$ and hence (2.8) will

have a unique continuous solution under the assumption $A_1 - A_3$. let us denote such a solution by

$$\mathrm{B}(t) = B^T(P_t), \quad t \in [0, T]$$

Or

$$B_1(t) = B_1^T(P_{1t}, P_{2t}), \quad B_2(t) = B_1^T(P_{1t}, P_{2t}), \ t \in [0, T]$$

If we supply this solution $(B_1, B_2)$ in (2.6) we find that (2.) will be satisfied if and only if $(P_{1t}, P_{2t})$ is a fixed point of the

operator

$$\Pi(P_t) = \{\Pi_1(P_{1t}, P_{2t}),\ \Pi_2(P_{1t}, P_{2t})\}$$

Where $\pi_1$ *and* $\pi_2$ respectively denote the right sides of (2.6) for *i =1,2* .

It is not difficult to see from the nature of $M_i$ , $L_i$ and the hypotheses on $f_i, b_i, \varphi_i$ that

$$B: C_T^2 \to C_T^2 \text{ and } \Pi C_T^2 \to C_T^2$$

(which guarantees the nonnegativity of $P_i$ and $B_i$). We will first show that

There exist positive constants $r^*, t^*$ such that $\pi$ is a mapping of a sphere

$$S_{t*}(\Phi, r^*) = \left\{ (P_{1t}, P_{2t}) \in C_{t*}^2 \mid \|P_t - \Phi_t\|_{t*} \le r^* \right\}$$

Into itself and also is a contraction where $\Phi_t$ is defined on [-$\tau$, $t^*$]

By the following:

$$\Phi_t = (\Phi_{1t}, \Phi_{2t});$$

(2.12)

$$\Phi_{it} = \begin{cases} \int_0^\infty \varphi_i(a, s)\, da \text{ for } s \in [-\tau, 0]. \\ \int_0^\infty \varphi_i(a, 0)\, da \text{ for } s \in [0, t^*]. \end{cases}$$

We note that similar to (2.11) one gets the estimates

$$|\Pi(P)(t) - \Phi_i(0)| \le \{B_i^* \Phi_i(0) \exp[B_i^* t] + \Phi_i(0) d_i^*\} t$$

Which together with (2.7) and (2.12) lead to

(2.13) $$\|\Pi(P_t) - \Phi_t\|_{t*} \le \left( \sum_{j=1}^2 \|\Phi_{jt}\|_{t*} [B_j^* \exp(B_j^* t) + d_i^*] \right) t^*$$

Thus if we choose,

$$r^* = T \sum_{j=1}^2 \|\Phi_{jt}\|_{t*} [B_j^* \exp(B_j^* T) + d_i^*]$$

Then $\Pi$ is a mapping of $S_{t*}(\Phi, r^*)$ into itself for all $t^* \le T$. we then will have to show that for small $t^* > 0$ , $\Pi$ is a contraction.

Let $P_t^{(1)}$ *and* $P_t^{(2)}$ be arbitrarily chosen in $S_{t*}(\Phi, r^*)$ and consider $|\Pi(P_t^{(1)}) - \Pi(P_t^{(2)})|$ for $t \le t^*$.

Corresponding to the chosen $P_t^{(1)}$ , $P_t^{(2)}$ we let $B^{(1)}(a) = B(P_T^{(1)})$ and $B^{(2)}(a) = B(P_T^{(2)})$ . Then from (2.6),

$$|\Phi(P^{(1)}(t)) - \Phi(P^{(2)}(t)| \leq \left| \int_0^t B_1^{(1)}(a)\left\{M_1\left(a,t,P_t^{(1)}\right) - M_1\left(a,t,P_t^{(2)}\right)\right\}da \right|$$

$$+ \left| \int_0^t [B_1^{(1)}(a) - B_2^{(2)}(a)]\, M_1(a,t,P_t^{(2)})da \right|$$

$$+ \left| \int_0^t \varphi_1(a,0)\left[L_1\left(a,t,P_t^{(1)}\right) - L_1\left(a,t,P_t^{(2)}\right)\right]da \right|$$

$$= J_1 + J_2 + J_3 \text{ (say)}$$

using the elementary inequality

$$|e^x - e^y| \leq |x-y| \ for \ x,y \leq 0$$

We derive from (2.5) that

$$\left| M_1(a,t,\ p_t^{(1)}) - M_1(a,t,\ p_t^{(2)}) \right|$$

$$\leq \int_o^{t-a} \left| f_1(s,\ p_t^{(1)}(a+s)) - f_1(s,\ p_t^{(2)}(a+s)) \right| ds$$

$$\leq d_1 \| p_t^{(1)} - p_t^{(2)} \|_{t^*} \text{(using } A_1\text{)}$$

*Now,using (2.10),*

$$\left| J_1 \right| \leq \int_o^a \beta_1^* \Phi_1(0)\, e^{\delta 1^a} \left(\int_o^{t-a} d_1 \| P_t^{(1)} - P_t^{(2)} \|_{t^*} ds\right) da$$

$$\leq d_1 \beta_1^* \Phi_1(0)\, t^2 \| P_t^{(1)} - P_t^{(2)} \|_{t^*} e^{\delta 1^t/2}$$

In a similar way

$$\left| J_3 \right| \leq \int_o^\infty \varphi(a,o)\left[\int_o^t \left| f_1(a+s, P_t^{(1)}(s) - f_1(a+s, P_t^{(2)}(s)) \right| ds\right] da$$

$$\leq d_1 \Phi_1(0) \| P_t^{(1)} - P_t^{(2)} \|_t$$

Now we have to estimate $\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right|$ in order to estimate $J_2$, from (2.8),

$$B_1^{(1)}(a) - B_1^{(2)}(a)$$

$$= \int_o^a \left[ B_1^{(1)}(\sigma) - B_1^{(2)}(\sigma) b_1 \left( a - \sigma, P_t^{(1)}(a) \right) M_1(\sigma, a, P_t^{(1)}(a) d\sigma + \int_o^a \left[ B_1^{(2)}(\sigma) \right. \right.$$

$$\text{X} \left[ b_1(a - \sigma, P_t^{(1)}(a) M_1(\sigma, a, P_t^{(1)}(a) - b_1 \left( a - \sigma, P_t^{(1)}(a) M_1 \left( \sigma, a, P_t^{(2)}(a) \right) \right] d\sigma \right.$$

$$+ \int_o^\infty \varphi(\sigma, o) \left[ b_1 \left( \sigma + a, P_t^{(1)}(a) \right) - b_1 \left( \sigma + a P_t^{(2)}(a) \right) \right], L_1 \left( \sigma, a, P_t^{(1)}(a) \right) d\sigma$$

$$+ \int_o^\infty \varphi(\sigma, o) b_1 \left( \sigma + a, P_t^{(2)}(a) \right) \left[ L_1 \left( \sigma, a, P_t^{(1)}(a) \right) - L_1 \left( \sigma, a, P_t^{(2)}(a) \right) \right] d\sigma$$

$$= R_1 + R_2 + R_3 + R_4 \text{ (say)}.$$

By the uniform Lipschitz continuity of $b_1$ in $A_2$ we get

$$\left| b_1 \left( a + \sigma, P_t^{(1)}(a) \right) - b_1 \left( \sigma + a, P_t^{(2)}(a) \right) \right| \leq B_1 \, \| P_t^{(1)} - P_t^{(2)} \| \text{ for } a \, \varepsilon [o, t^*]$$

And hence we have

$$|R_3(a) + R_4(a)| \leq \beta_1 \Phi_1(0) || P_t^{(1)} - P_t^{(2)} ||_{t^*} + \Phi_1(0) \beta_1 d_1 || P_t^{(1)} - P_t^{(2)} ||_{t^*}^{*a},$$

$$|R_2(a)| = [\int_0^a [B_1^{(2)}(\sigma) \{ b_1( a-\sigma, P_t^{(1)}(a)) - b_1(a-\sigma, P_t^{(2)}(a)) \} M_1(\sigma, a, P_t^{(1)}(a))$$

$$+ b_1(a, -\sigma, P_t^{(2)}(a)) \{ M_1(\sigma, a, P_t^{(1)}(a)) - M_1(a-\sigma, P_t^{(2)}(a)) \}] d\sigma]$$

$$+ b_1 \left( a - 0, P_t^{(2)}(a) \right) \{ M_1(\sigma, a, P_t^{(1)}(a)) - M_1(a-\sigma, P_t^{(2)}(a)) \}] d\sigma]$$

$$\leq \int_0^a [\beta_1^* \Phi_1(0) \, e^{\delta 1^a} \beta_1 \| P_t^{(1)} - P_t^{(2)} \|_{t^*} + \beta_1 d_1 \| P_t^{(1)} - P_t^{(2)} \|_{t^*} (a-o)] do$$

$$\leq \beta_1^* \Phi_1(0) \, e^{\delta 1^a} \| P_t^{(1)} - P_t^{(2)} \|_{t^*} \{ \beta_1 a + d_1 a^2 / 2 \}$$

Thus

$\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right|$

$$\leq \int_o^a \left| B_1^{(1)}(\sigma)\text{-}B_1^{(2)}(\sigma) \right| B_1^{*d}\sigma$$

$$+ \Phi_1(0) /\!/ P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} \left\{ B_1 + B_2\, d_1\, a + B_1^{*} \left( B_2 a + d_1 a^2/2 \right) e^{\delta 1^a} \right\}$$

$$\leq \beta_1^* \int_o^a \left| B_1^{(1)}(\sigma) - B_1^{(2)}(\sigma) \right| \mathrm{do} + \left( r^* + /\!/ \Phi_t /\!/_{t*} \right) \| P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} g_1^{(a)}$$

Where $g_1(a) \rightarrow \beta_1$ as $a \rightarrow 0$. By Gronwall's inequality it will follow that

$$\left| B_1^{(1)}(a)\text{-}B_1^{(2)}(a) \right| \leq \left( r^* + /\!/ \Phi_t /\!/_{t*} \right) /\!/ P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} \{ g_1(a) + 0(a) \}$$

Using this estimate for $R_1(a)$ we get

$$|R_2(a)| \leq \left| \int_0^a [B_1^{(2)}(\sigma)\, B_1^{(2)}(\sigma),]\, b_1(a)) - b_1(a\text{-}\sigma, P_t^{(1)}(a))\, M_1(\sigma, a, P_t^{(1)}(a)) do \right|$$

$$\left| B_1^{*} \left( r^* + /\!/ \Phi_t /\!/_{t*} \right) /\!/ P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} 0(a) as\ a \rightarrow 0 \right.$$

*Thus*

$$\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right| \leq (r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ P_t^{(1)} - P_t^{(2)} /\!/_t o(1) as \quad a \rightarrow o$$

And hence

$$\left| \pi_1(P_t^{(1)} - \pi_1\left(P_t^{(2)}\right) \right| \leq ((r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ P_t^{(1)} - P_t^{(2)} /\!/_t o(t) as \quad t \rightarrow o$$

Which together with a similar bound for $\left| \pi_2(P_t^{(1)} - \pi_2\left(P_t^{(2)}\right) \right|$ leads to

$$\| \pi\,(P_t^{(1)} - \pi\left(P_t^{(2)}\right) /\!/ \leq (r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ P_t^{(1)} - P_t^{(2)} /\!/_{t*} o(t^*)$$

As $t^* \rightarrow o$ showing that if $t^*$ is small enough, the mapping $\pi$ on

$s_{t^*}(\Phi, r^*)$ is a contraction. Hence the system (2.6)-(2,8) has a unique solution for $o \leq t \leq t^* for\ t^*$ sufficiently small and positive

Now if [0,T] is any finite interval on $R^+$, we can consider [0,T] as a finite union of intervals of length less than or equal to $t^*$ Since the estimate (2.11)is valid for any finite interval we can extend the solution of (2,6)-(2.8) form $[0, t^*]$ to [0,T] and this completes the proof,

3. Stationary age distributions and their linear stability we will now establish the existence of time independent solution

$P_t^*, P_t^*$ Of (1.1)-(1.6): such solutions satisfy

(3.1) $\frac{dp\,P_2^*(a)}{da}+f_i(\mathrm{a}, P_1^*, P_2^*)\,P_i^*(a) = 0,\ \ i = 1,2, a > 0,$

(3.2) $P_i^* = \int_o^\infty P_i^*(a)da, i = 1,\ 2,$

(3,3) $P_i^*=\int_o^\infty b_i(a,\ P_1^*,\ P_2^*)P_i^*(a)da,\ i = 1.\ \ 2$

All solutions of (3.1)are of the from

(3.4) $P_i^*(a) = P_i^*(o)\exp\left[-\int_o^\infty f_i(s, P_1^*, P_2^*)ds\right], i = 1,\quad 2, a \geq 0,$

Which together with (3.3) lead to

$1=\int_o^\infty b_1(a, P_1^*, P_2^*)\exp\left[-\int_o^a f_1(s, P_1^*, P_2^*)ds\right]da,$

(3.5) $1=\int_o^\infty b_2(a, P_1^*, P_2^*)\exp\left[-\int_o^a f_2(s, P_1^*, P_2^*)ds\right]da$

Thus the existence of stationary solutions of (1.1)-(1.6) reduces to the existence of a pair of positive constants $P_1^*, P_2^*$ satisfying (3.5). We can now prove the following,

THEOREM 3. Assume that $f_1$ and $b_i$ satisfy ( $A_1$ ) and ( $A_2$ )and let

(3.6) $F_{i.}\ (P_1, P_2) = \int_o^\infty b_i(a, P_1, P_2)\exp\left[-\int_o^\infty f_i(s, P_1, P_2)ds\right]da,$

$i$ =1, 2, $P_1 \geq 0$, $P_2 \geq 0$,

**And suppose the following hold:**

(3.7) $A_4$. $F_{i.}\ (0,\ 0) > 1,\ i = 1,2$:

$A_5$. *there exists a positive constant c for which*

$F_1$(C,0)< $F_{2.}$(C,0),

(3.8) $F_1$(0.C)< $F_{2.}$(0,C),

Then there exists a unique pair $(P_1^*, P_2^*)$ of real numbers $P_1^* > o$,

$P_2^* > 0$ such that

(3,9) $F_1(P_1^*, P_2^*) = 1 = F_2(P_1^*, P_2^*)$

And the unique nonegative solution of (3.1)-(3.3) is given by

(3.10) $P_2^*(a) = [P_2^* \exp[-f_i(s, P_1^*, P_2^*)ds])/(\int_o^\infty \exp[-\int_o^a f_i(s, P_1^*, P_2^*)ds]da)$,

$i = 1, 2, \quad a \geq 0,$

Proof Consider the elements of the surfaces defined by

$$Z = F_1(P_1, P_2), z = F_2(P_1, P_2), \ P_1 \geq 0, P_2 \geq 0,$$

In the ($P_1$,$P_2$,z) space. By hypothesis ($A_1$) we have $\partial F_i / \partial P_j < 0$ *and* , by ($A_2$), $\partial b_i / \partial P_j \leq 0$ for all

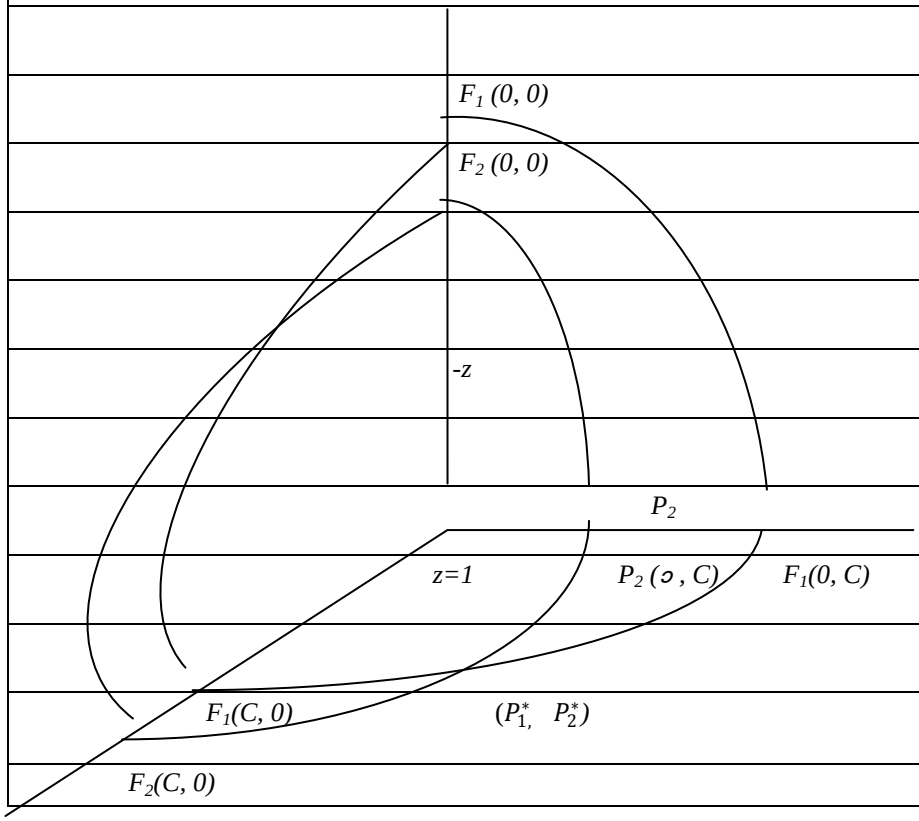
$$P_1 \geq 0, P_2 \geq 0, \quad i = 1, 2,$$

$j = 1, 2,$ Hence by ($A_4$) the intersections of these surfaces with the plane z=1 define two curves on the z=1 plane connecting the I ines {z=1, $P_1 \geq 0, P_2 \geq 0,$} and {z= $1, P_2 \geq 0, P_1 \geq 0$, These two curves are defined by

$$F_1(P_1, P_2 = 1 \ and \ F_2(P_1, P_2) = 1, \ F_1(P_1, P_2 \geq 0,$$

Which by ($A_s$) will intersect at a unique point say $((P_1^*, P_2^*)$ with $(P_1^* > 0, P_2^* > 0$ (see figure). The uniquess of the point $((P_1^*, P_2^*)$ *is a* consequence of the smoothness of the surfaces $z = F_1(P_1, P_2)$ and $z = F_2(P_1, P_2)$

Following is a hypothetical case for Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates attributable or ascribable in the transactions and reactionary potentialities.



(The figure corresponds to the case $F_1(0,0) > F_2(0,0) > 0$. The other possibility $1 < F_1(0,0)$ is treated similarly) It will now

Follow that for such a pair $(p^*_1, p^*_2)$ the unique nonnegative solution of (3.1) - (3.3) is given by (3.10).

Now to examine the linear stability of the stationary age distributions $p^*_1, p^*_2$, we let

(3.11)
$$\begin{cases} p_i(a,t) = p^*_i(a) + u_i(a,t) \\ p_i(t) = p^*_i + p_i(t), \quad i = 1, 2, \quad a > 0, \quad t > 0, \\ p_i(t) = \int_0^\infty u_i(a,t)da, \end{cases}$$

In (1.1)-(1.5) and derive the following variational system after neglecting the nonlinear terms in the perturbations $\boldsymbol{u}_i$ and $\boldsymbol{p}_i$ $(i = 1, 2,)$ (such a procedure of linear stability analysis can be justified As has been done in Gurtin and MacCamy [4].)

(3.12)
$$\begin{cases} Du_i(a,t) = -f_i\left(a, p^*_1, p^*_2\right) u_i(a,t) - p^*_i(a)\left\{\sum_{i=1}^{2} \frac{\partial f_i}{\partial p^*_i} p_i(t)\right\}, \\ u_i(0,t) = \int_0^\infty b_i\left(a, p^*_1, p^*_2\right) u_i(a,t)da + \int_0^\infty p^*_i(a)\left\{\sum_{i=1}^{2} \frac{\partial b_i}{\partial p^*_i} p_i(t)\right\} da \end{cases}$$

Where $i = 1,2$ and (3.13)
$$p_{11}(t) = p_1(t), \quad p_{12}(t) = p_2(t-\tau), \quad p_{21}(t) = p_1(t-\tau), \quad p_{22}(t) = p_2(t)$$

To consider the asymptotic bahaviour as $t \to \infty$ as of solutions of (3.12) we let

(3.14)
$$\begin{cases} u_i(a,t) = u_i(a) \exp[\lambda t], \quad u_i(a) \to 0 \text{ as } a \to \infty, \\ p_i(t) = p^*_i \exp[\lambda t], \\ \qquad\qquad i = 1, 2, \quad a > 0, \quad t > 0, \\ p^*_i = \int_0^\infty u_i(a)da, \end{cases}$$

Using (3.14) in (3.12) - (3.13),

$$\frac{du_i}{da} + \lambda u_i = -f_i\left(a, p^*_1, p^*_2\right) u_i(a) - p^*_i(a)\left\{\sum_{j=1}^{2} \frac{\partial f_i}{\partial p^*_j} p^*_i\right\},$$

(3.15)

$$\mu_i(0) = \int_0^\infty b_i\left(a, p_1^*, p_2^*\right)\mu_i(a)da + \int_0^\infty p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial b_i}{\partial p_j^*}\ p_i^*\right\}da,$$

Where $i = 1,\ 2,$ and

$$p_{11}^* = p_1^*, \quad p_{12}^* = p_2^* \exp\{-\lambda\tau\},$$

$$p_{21}^* = p_1^*, \exp\{-\lambda\tau\}, \ p_{22}^* = p_2^*.$$

If we integrate the first of (3.15) and use the fact that $\mu_i(a) \to 0$ as $a \to \infty$ we can derive that $p_1^*$ and $p_2^*$ are governed by

(3.16) $$\lambda p_1^* - \int_0^\infty p_1^*(a)\left[\left(\frac{\partial b_1}{\partial p_j^*} - \frac{\partial f_1}{\partial p_1^*}\right)p_1^* + \left(\frac{\partial b_1}{\partial p_2^*} - \frac{\partial f_1}{\partial p_2^*}\right)p_2^* e^{-\lambda\tau}\right]da$$

$$= \int_0^\infty \left[b_1\left(a, p_1^*, p_2^*\right) - f_1\left(a, p_1^*, p_2^*\right)\right]\mu_i(a)da,$$

(3.17) $$\lambda p_2^* - \int_0^\infty p_2(a)\left[\left(\frac{\partial b_2}{\partial p_1^*} - \frac{\partial f_2}{\partial p_1^*}\right)p_1^* e^{-\lambda\tau} + \left(\frac{\partial b_2}{\partial p_2^*} - \frac{\partial f_2}{\partial p_2^*}\right)p_2^*\right]da$$

$$= \int_0^\infty \left[b_2\left(a, p_1^*, p_2^*\right) - f_2\left(a, p_1^*, p_2^*\right)\right]\mu_2(a)da,$$

From the first of (3.15) we get

(3.18) $$\mu_i(a) = \left\{\mu_i(0)p_i^*(a)/p_i^*(0)\exp[\lambda a]\right\}$$

$$- p_i^*(a)\exp[-\lambda a]\int_0^a \left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right\}e^{\lambda s}ds, \quad i = 1, 2,$$

Which by the second of (3.15) becomes,

(3.19) $$\mu_i(a)\left[p_i^*(a)/p_i^*(0)c_i\exp(\lambda a)\right]\left[\int_0^a p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right\}da\right.$$

$$\left. - \int_0^\infty b_i\left(a, p_1^*, p_2^*\right)p_i^*(a)\,e^{-\lambda a}\left\{\int_0^a\left(\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right)e^{\lambda s}ds\right\}da\right]$$

$$- p_i^*(a)\,e^{-\lambda a}\int_0^a\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right\}e^{\lambda s}ds, i = 1, 2,$$

where

(3.20) $$c_i = 11 - \frac{1}{p_i^*(0)} \int_0^\infty b_i\left(a, p_1^*, p_2^*\right) p_i^*(a) e^{-\lambda a} da, \quad i = 1, 2,$$

Now supplying $\Phi_1(a)$ and $\Phi_2(a)$ from (3.19) in (3.16)-(3.17) we derive that

$$\left(\lambda - A_{11}\right) p_1^* - A_{12} p_2^* \exp\left(-\lambda\tau\right) = g_{11}(\lambda) p_1^* + g_{12}(\lambda) p_2^* \exp(-\lambda\tau),$$

(3.21)

$$-A_{21} p_1^* \exp\left(-\lambda\tau\right) + \left(\lambda - A_{22}\right) p_2^* = g_{21}(\lambda) p_1^* \exp\left(-\lambda\tau\right) + g_{22}(\lambda) p_2^*$$

Where

(3.22) $$A_{i\Phi} = \int_0^\infty p_i^*(a) \left[\frac{\partial b_i}{\partial p_\Phi^*} - \frac{\partial f_i}{\partial p_\Phi^*}\right] da,$$

(3.23) $$g_{i\Phi}(\lambda) = \left[p_i^*(0) c_i\right]^{-1} \left\{\int_0^\infty \sigma_i(a) p_i^*(a) e^{-\lambda a} da\right\} \Phi \left\{\int_0^\infty \left[p_i^*(a) \frac{\partial b_i}{\partial p_\Phi^*} - b_i\left(a, p_1^*, p_2^*\right) p_i^*(a) e^{-\lambda a} \left(\int_0^a \frac{\partial f_i}{\partial p_\Phi^*} e^{\lambda s} ds\right)\right] da\right\}$$

$$- \int_0^\infty \sigma_i(a) p_i^*(a) e^{-\lambda a} \left(\int_0^a \frac{\partial f_i}{\partial p_\Phi^*} e^{\lambda s} ds\right) da, \quad i, \Phi = 1, 2,$$

(3.24) $$\sigma_i(a) = b_i\left(a, p_1^*, p_2^*\right) - f_i\left(a, p_1^*, p_2^*\right), \quad i, = 1, 2.$$

To solve the linear perturbational system (3.12) it is enough to find the constants $\lambda, p_1^*, p_2^*$ satisfying (3.21) since one can then use (3.19) to find the perturbations from (3.14). A nontrivial solution $\left(p_1^*, p_2^*\right)$ for (3.21) will exist if and only if $\lambda$ is a root of the equation $et \begin{bmatrix} \lambda - \left(A_{11} + g_{11}(\lambda)\right) & -\left(A_{12} + g_{12}\right) e^{-\lambda\tau} \\ -\left(A_{21} + g_{21}\right) e^{-\lambda\tau} & \lambda - \left(A_{22} + g_{22}(\lambda)\right) \end{bmatrix} = 0$

Or equivalently

(3.25) $$\lambda^2 - \lambda\left(A_{11} + A_{22}\right) + A_{11}A_{22} - A_{12}A_{21} \exp\left[-2\lambda\tau\right] = \lambda s_1(\lambda) + s_2(\lambda, \tau)$$

Where $S_1(\lambda) = g_{11}(\lambda) + g_{22}(\lambda)$ ,

$S_2(\lambda, \tau) = \left[A_{21} g_{12}(\lambda) + A_{12} g_{21}(\lambda) + g_{12}(\lambda) g_{21}(\lambda)\right] \exp\left[-2\lambda\tau\right]$

$$-A_{11} g_{22}(\lambda) - A_{22} g_{11}(\lambda) - g_{11}(\lambda) g_{22}(\lambda)$$

The following result establishes the linear local asymptotic Stability of the stationary age distributions $p_1^*(a)$ and $p_2^*(a)$

THEOREM 4 Assume the following:

$A_6.\ 0 < \beta_i^* < d_i^*\ (i,\ =1,\ 2.)\ \ (see(A_1)\ and\ (A_2));$

$A_7.\ \{A_{11} + A_{22} + \| S_1(0) \|\}^2 < 4\{(A_{11}A_{22} - A_{12}A_{21}) - \| S_2(0) \|\},$

$$\| S_1(0) \| = |g_{11}(0)| + |g_{22}(0)|,$$

$$\| S_2(0) \| = |A_{11}\, g_{12}(0)| + |A_{12}\, g_{21}(0)| + |g_{12}(0)\, g_{21}(0)|$$

$$+ |A_{11}\, g_{22}(0)| + |A_{22}\, g_{11}(0)| + |g_{11}(0)\, g_{22}(0)|$$

Then all the roots of (3.25) have negative real parts and hence the

perturbations $u_i\ (a,\ t) = \phi(a) \exp[\lambda t] \to 0\ as\ \ t \to \infty$ for $i = 1, 2\ and\ a \geq 0$

Proof . Define $F_1(\lambda, \tau)$ and $F_2(\lambda, \tau)$ as follows:

$F_1(\lambda, \tau) = \lambda^2 - \lambda(A_{11} + A_{22}) + A_{11}A_{22} - A_{12}A_{21} \exp[-2\lambda\tau];$

$F_2(\lambda, \tau) = \lambda S_1(\lambda) + S_2(\lambda, \tau),$

Using $(A_6)$ one can show that for $\lambda$ with re $(\lambda) \geq 0$, $c_1$ and $c_2$ are positive and bounded away from zero, Now since $A_{ij} > 0$ and $\tau \geq 0$ we find from the nature of the dependence of $g_{ij}$ on $\lambda$ that for all $\lambda\ \ \mu \pm iw$ with $\mu \geq 0$ and $w \geq 0$ we have

(3.26) $$\left| F_1(\lambda, \tau) - P_2(\lambda, \tau) \right|$$

$$\geq \left| F_1(\lambda, 0) | - F_2(\lambda, \tau) \right|$$

$$\geq |\lambda|^2 - |\lambda| \{A_{11} + A_{22} + \| S_1(0) \|\} + \{A_{11}A_{22} - A_{12}A_{21} - \| S_2(0, \tau) \|\}$$

$$\geq \left[ (|\lambda| - P)^2 + (q) \right]^2$$

Where

$p = \{A_{11} + A_{22} + \| S_1(0) \|\}/2,$

$q = \{4 | (A_{11}A_{22} - A_{12}A_{21}) - \| S_2(0, T) \| | - | A_{11} + A_{22} + \| S_1(0) \| |^2\}/2,$

It will now follow, from (A) and (3.26), $F_1(\lambda, \tau) = F_2(\lambda, \tau)$ connot have roots with zero or positive real parts and this completes the proof.

In conclusion we remark that the condition (3.7) means that the net creation rates of each system is greater than utility; a condition of this type is well known in age wise classification of particles, atoms, molecules, cosmological phenomenon. (3.8) can be interpreted to mean that the intraspecific inhibition is higher than that of the

interspecific interaction. The conditions of Theorem 4 are analytical and do not lend themselves for any worthwhile interpretation in terms of the model parameters. So also an interpretation can be recorded as a remark in respect of destruction rates of particles, properties, elements, or projections.

Note: The model is applicable to each and every module. As stated in the foregoing, there are possibilities of situation where an account itself be a receiver and it is also the one to which debit is raises. Such examples are legion in all the fields of finance, economics, political science and physics and chemistry, and of course mathematics. Classification based on time is the preeminently and predominantly suited premise for such model. We repeat that the model is to be read with each module We will expatiate the importance and implicational ramifications as we proceed with the work. Assortment-blend, collection-combinations, combo-gallimaufry, goulash-hash, hodgepodge-medley, mishmash-mixture, motley-pastiche patchwork-salmagundi, soup-stew, mixed bag-mélange, olio- olio podrida, assortment variety, Array choice, collection combination, combo diversity, hodgepodge jumble, kind medley, miscellany mishmash, mixed bag mixture, mélange potpourri, selection sort, brew concoction, beverage blend, broth compound, distillation drink, fermentation hash, hodgepodge infusion, instillation liquor, mélange miscellany, mishmash mixture, mélange potpourri, borscht bouillon, bowl brew, chowder concoction, consommé decoction, gumbo hodge-podge, olio porridge, potage potpourri, pottage puree, splash stock, vichyssoise hodgepodge, mixture- mess, collection combination, goulash hash, jumble medley, miscellany mishmash, mixed bag mélange, olio patchwork, potpourri salmagundi, jumble clutter, disarrangement disarray, disorder farrago, gallimaufry garbage, goulash hash, litter medley, mess miscellany, mishmash mixture, muddle mélange, olio pastiche, patchwork potpourri, salmagundi scramble, shuffle snarl, tangle tumble, medley miscellany, assortment brew, collection combo, conglomeration farrago, hodgepodge jumble, melee mingling, mishmash mixed bag, mixture mélange, pasticcio pastiche, patchwork potpourri, salmagundi variety, mélange mixture, assortment combo, farrago gallimaufry, pasticcio pastiche, patchwork potpourri, salmagundi soup, stew miscellany, accumulation aggregation, anthology assortment, brew collectanea, combination combo, compilation conglomeration, accumulation diversity, farrago gallimaufry, melee mishmash, mix mixed bag, mixture muddle, mélange odds and ends, olio pasticcio, pastiche patchwork, potpourri salad, salmagundi smorgasbord, stew variety, mishmash jumble, medley mess, miscellany mixed bag, mixture mélange, olio patchwork, potpourri salmagundi, mixed bag array, assortment choice, collection combination, combo conglomerate, diversity hodgepodge, jumble medley, mélange miscellany, mishmash mixture, montage mélange, odds and ends, patchwork potpourri, smorgasbord sundries, mix/mixture assortment, combination, admixture adulteration, alloy amalgam, amalgamation assimilation, association batter, blend brew, combine combo, commixture composite, compound concoction, confection conglomeration, cross crossing dough fusion, goulash grab bag, hodgepodge hybrid, hybridization incorporation, infiltration interfusion, jumble mash, medley merger, mingling miscellany, mishmash mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, odds and ends assortment, bits and pieces, debris etcetera, hodgepodge jumble, knickknacks leavings, leftovers litter, medley melee, miscellaneous paraphernalia, miscellany mishmash, mixed bag motley, mélange notions, novelties oddments, olio particles, patchwork potpourri, rags refuse, remainder remains, remnants rest, rummage scraps, sundries sundry items, compilation mishmash, paste-up patchwork, potpourri reappropriation, reproduction synthesis, conglomeration cross, crossing dough, fusion goulash, interfusion jumble, mash medley, merger mingling, miscellany mishmash, mix mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, colluviums gallimaufry, hash mixture, pasticcio pastiche, conglomeration cross section, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multiformity multiplicity, mélange potpourri, range shift, variegation variousness, divergencies diverseness, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multifariousness multiformity,

mélange potpourri, range shift, soup stew, variance variation, variegation variousness, oddment odds and ends, assortment bits, polymorphism assortment, change collection, combo conglomeration, cross section departure, discrepancy disparateness, divergencies diverseness, diversification diversity, fluctuation heterogeneity, heterogeneousness incongruity, intermixture many-sidedness, medley miscellaneousness, miscellany mishmash, mixed bag mixture, notions novelties, oddment oddments, olio particles, patchwork potpourri of the particles, interactions, accounts, transactions in the cosmos. The wisdom and understanding of our own divine essence is like a beaconing begorra, bureaucratic burdock, befuddling binary and digital digamma, that invites us inward toward ever deeper connection with all around us. Yet rather than look within, most in this world prefer to look outward for divine wisdom, thus sustaining a dependence upon a vast hierarchy that stretches between the individual and the Divine. In all our wanderings away from the Divine, humankind has obscured, obfuscated, obliterated, adjured, abnegated and abjured its most compelling features through a persistent belief systems in limitations arising out of the controlling mechanisms of this hierarchy of sub ordinates, superordinates and coordinates. Has my consciousness changed? Why that doubt? Twist and shake, ponderous novelty themes, pathos, poignancy in writing in different styles, accord-coetaneousness, coevality-coincidence, concurrence-conformity, conjunction-contemporaneousness, harmony, order, peace, simultaneousness- synchronicity exciting optimism of youth, power obmbudsman,pulsations and throbbing of pop, hard luck stories of people, modernity and mischief, grimalkins of rurality, wobbling boats of lovers, saxophone in mouths all point to newer value orientations. The Divine dances and the “Grand design “is perpetrated outside of the confines of any hierarchical structures, be it feudal, organizational. It is complete within itself, and has a singular purpose of demonstrating the collective potential of all life within the universe. It is the archetype of perfection. It is the standard bearer of each soul's innate design and ultimate destiny. The essence of the Divine is far beyond mental conception, yet humanity's tendency is to resort to the limiting language of the hierarchical paradigm acceding- accession, accommodation- accord, accordance, acknowledging, adjustment-affiliation, affinity-alliance, amity-approving, arbitration-arrangement, assenting, authorizing, bargaining-compatibility, compliance--compromise, concert-concession, conformity-congruity, consistency-correspondence, endorsing, granting-harmony in conformity to pen chance ,predilections of the protagonist to define and understand it. Searching for Connection and Wholeness calls for not lumping elements like me made of clay feet and weak knees, owing allegiance to the fringe, without courante and courage of conviction and will for vindication, visual restlessness and exaggerated narratives of self to its own, exciting potentialities available, terrestrial passion, and skew of scams all bear an infallible demonstration for new order of consciousness. Theorists have posited a virtual smorgasbord of objects to account for dark matter, although many of them have fallen prey to observational constraints. As leading possible candidates for baryonic dark matter, there are black holes (large and small), brown dwarfs (stars too cold and faint to radiate), sun-size MACHOs, cold gas, dark galaxies and dark clusters, to name only a few. The range of particles that could constitute nonbaryonic dark matter is limited only slightly by theorists' imaginations. The particles include photinos, neutrinos, gravitinos, axions and magnetic monopoles, among many others. Of these, researchers have detected only neutrinos--and whether neutrinos have any mass remains unknown. Experiments are under way to detect other exotic particles. If they exist, and if one has a mass in the correct range, then that particle might pervade the universe and constitute dark matter. But these are very large "ifs." To a great extent, the details of the evolution of galaxies and clusters depend on properties of dark matter. Without knowing those properties, it is difficult to explain how galaxies evolved into the structures observed today. As knowledge of the early universe deepens, I remain optimistic that we will soon know much more about both galaxy formation and dark matter. What we fail to see with our eyes, or detectors, we can occasionally see with our minds, aided by computer graphics. Computers now play a key role in the search for dark matter. Historically, astronomers have focused on observations; now the field has evolved into an experimental science. Today's astronomical experimenters sit neither at lab benches nor at telescopes but at computer terminals. They scrutinize cosmic simulations in which tens of thousands of points, representing stars, gas and dark matter; interact gravitationally over a galaxy's lifetime. A cosmologist can tweak a simulation by adjusting the parameters of dark matter and then watch what happens as virtual galaxies evolve in isolation or in a more realistic, crowded universe.

Old consciousness will not do. Newer one has to be imbibed and cultivated. When people are unaware of their inner wisdom and wholeness, they tend to search for order and security outside themselves. That is my personal experience. There is always search for sanctity of emotion, aesthetic expression, pursuit of harmony, unraveling realms, syncopated ending and stillness of mind and shiver of heart at the mirror broken somewhere are the strains of newer wisdom if it can be called one. Same questions shall be asked and same shall be answer in the similar vein. Uncertain of their place within the hierarchical order of the world, they define themselves based upon their insecurities. Individuals thus become only pieces of their wholeness and like shards of glass from a beautiful vase; they bear little resemblance to their aggregate beauty. Within the hierarchy, many in high positions of power have taken advantage of our collective insecurities in attempting to guide the development of all humankind. They have obscured the direct connection between the individual and the Divine through a variety of means designed to intercede between our inner essence and our divine source. Each individual can come to know them to be free of all forms of institutional and organisational and societal hierarchical control. This is not to imply that we should not trust others or join together in bonds of friendship and community. It is simply a reminder that relative truth is constantly shifting in the hands of those who desire to control. And even when the motive for their controlling behaviors may be of good will, it is still a form of control. When the revealers of "truth" within the hierarchy withhold and suppress information, they are usually positioning themselves to maintain or acquire more power rather than to disseminate empowerment to all. The desire for connection and wholeness is a fuel that drives us to seek out and explore the hierarchical paradigm. This inner longing provides us with the motivation to seek help and guidance from a specific group within the hierarchy, and in so doing, cultivate a sense of belonging and connection. Furthermore, the hierarchical paradigm is a stage whereby we develop a sense of connection to some grand, encompassing vision. This is why the hierarchy nurtures prophets who point toward a greater vision. Spiritual leaders are able to peer deeply beneath the surface reality of life and experience how intricately connected every life form is, and how the composite of all life is intelligent far beyond measure. These visionary leaders can thus interpret reality through their personal abilities to perceive and express life's dimensional depth and limitless intelligence. Yet no one is able to articulate life's full dimensional depth and breadth with the tools of language. They can only, at best, describe their interpretation or their impressions. In actuality, all of us are able to peer beneath the surface reality of life and perceive a unique vision of the universe. We require only time and intention to develop our own interpretations. And this is precisely what many great spiritual leaders have taught. Life's deeper essence is not an absolute to be experienced by the chosen few, but an evolving, dynamic intelligence that wears as many faces as there are life forms. No group or people have the exclusive portal into the universe by which the Divine expresses itself in all its majesty. The portal is open and available to all, because the Divine is within all things. Those recognized as great prophets each produced a vision of the universe beyond what was currently defined by the hierarchy. Because their interpretations were articulated with authority and depth of insight, they became a target of debate among various groups in the hierarchy. This debate then created a polarity of belief. A sympathetic constituency emerged to defend and embellish their leader's interpretation, while established groups held it in contempt of previously held beliefs. Invariably, the leader's vision became confined and shaped into dogma by followers who desired to create a new religion or sect. Thus, this infusion of fresh insight quietly receded into the hands of the hierarchy, where its deeper meaning was obscured by the very fact that it was incorporated into a rigid structure that both protected and promoted it. A new paradigm is emerging that promotes a clear connection of individual consciousness to the compelling features of the Divine without the intervention of a hierarchy. This is when the fables and myths of history step into the light and become known as they were originally intended. This is the time when language will be transformed into a new form of communication that breaks down all barriers of control. Personal transformation, through the awakening of inner wisdom and understanding, is the pathway into wholeness. The transformation paradigm is initiated simply by the recognition that rather than the dependency-inducing ways of the hierarchical paradigm, there are accelerated, independent pathways that bypass the hierarchy and lead to self-mastery. These new pathways lead to the divine wisdom and understanding that is present within all of us. This wisdom can be accessed through the practice of three principles of transformation: seeing the Divine in

all, nurturance of life, and gratitude. The application of these life principles disengages individuals from the controlling elements of the hierarchy, thus initiating the transformation experience. As there are relative truths, there are relative freedoms. As individuals evolve through the hierarchical process, an ever-increasing sense of freedom is gained, yet external forces continue to exert control through limiting language, confining belief systems, and more. These controlling influences lead to anatomical architecture, body, build, cagey carcass, construction, enclosure, fabric, flounce, form, framework, fringe, groundwork, hem, mount, mounting outline, physique, scaffold, scaffolding scheme, setting shell, stage, structure, support, system, trim, trimming, truss- valance continued reliance upon the hierarchy as it unceasingly attempts to impose a sense of inequality between us. The underlying equation of the hierarchical paradigm is: individual + hierarchy = God connection. In the case of personal transformation it is: individual + inner wisdom and understanding divine equality with all. The time has come to integrate the dominant paradigm of the hierarchy with the liberating transformation paradigm. That shall redeem all promises made to oneself. There shall be a Matrix force or balancing force that shall remain UN satisfies until conservativeness is reached. This integration occurs naturally once we have fully explored the two paradigms and develop a synthesis paradigm whereby transformation is attained by searching for connection and wholeness through our own inner wisdom and understanding. It is this combination of release from dependency on the hierarchy and transformation from within that initiates the synthesis paradigm. Once we take responsibility for our own transformation and integration, it does not mean that the hierarchy is to be shunned or avoided. The hierarchy is quite benign as a manipulative force and merely represents one important stage in the journey toward wholeness. What is being set in motion now is the initial preparation for these paradigm shifts. More specifically, these paradigms will be simultaneously played out over the coming years. As always, it will be the choice of each one of us as to which paradigm we embrace in our journey. All of the highest imaginings of humanity are yet unaware of our deepest foundation. We have sought the upper reaches of the building, yet remain unaware of the foundation's design. It is here, at the very core of existence that the Divine is bursting forth with its creative energy while simultaneously reintegrating with its invitation to wholeness. It is here that equality is realized, not in the lofty places of relative truth lodged in the hierarchy, but rather in the deepest part of the foundational plan of life's original source and ultimate destiny. The origin and destiny of existence is the tone of equality in life. Listen for this tone – this frequency of vibration – and follow it back into the very foundation from whence all things arise and return. Consider these words as symbols only. Feel the truth that stands behind these words, and tap into this empowering energy force that reaches out for you. Know it as a tone or vibration – a resonance that waits for you around every corner in which your life will turn. It is a beacon of the Divine gathering itself into the form of language in order to usher you to a place from which you can experience the formless tone of equality – the bypass of limitation. It is the primal language of our divine source that bestows to you the freedom to generate your own deepest beauty in the expression of your highest truth.(Reference: Wikipedia)

## Gopalasamy Model For Procrastinated Conducive, Felicitous Counter Sallies In Two Specified Systems

Abstract

Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of decay rate of perturbations. Sole intention, main objective and primary concern of

the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are counted in the same vein as that of predator prey. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 7. Preliminary remarks

Almost all the equations of population ecology have been derived starting from the following simple format (Note this can be used for the development of the creed and type of systemic types in case of cosmology and astrophysics and fiancé and econometrics. To give an example of finance, we can think of a account of the company giving rise to various accounts of its employees, and that is the contribution of the single account for the development of the total system types of the given configuration). It may be account, particle, atom, molecule, proposition, problem or any such thing. We shall delineate and disseminate information and particulars as and when the model shall be applied to the problem in question.

$$\frac{dN(t)}{dt} = \text{an individual's contribution}$$

*to population growth* $N(t)$

(1.1)

where $N(t)$ denotes the size of a population of a single species (or biomass) with overlapping generations at time $t$, and $dN(t)/dt$ is the rate of change of the population size at time $t$. Subsequently one makes an assumption regarding the factor in the brackets in (1.1); in particular if one assumes that an individual entity's contribution to the population growth is denoted by a function say, f(t, N), defined suitably for all t>0, N then one obtains from (1.1) the so called Kolmogorov formulation in the form

$$\frac{dN(t)}{dt} = f(t, N(t))\, N(t).$$

(1.2)

Various choices of $f$ together with some ecologically plausible assumptions such as the temporal constancy of the environment and density dependent effects in $f$ lead to various equations. For instance if $f(t, N) \equiv r$ (a positive constant) one obtains the Malthusian formulation in the form

$$\frac{dN(t)}{dt} = rN(t)$$

And if one assumes $f(t, N) \equiv r - (r/K) N$ for some positive constants r and K one gets the familiar logistic equation of population biology

$$\frac{dN(t)}{dt} = rN(t) \left\{ 1 - \frac{N(t)}{K} \right\}$$

In order to device models with time delayed effects, such effects are introduced in the average growth rate $F(t, N)$ of (1.2); for instance if $f(t, N) \equiv r - (r/K) N(t - \tau)$ for some constant $\tau > 0$ we obtain Hutchinson's[19]delay loistic equation

$$\left\{ \frac{dN(t)}{dt} = rN(t) \left\{ 1 - \frac{N(t-\tau)}{K} \right\} \right. .$$

(1.3)

If for instance

$$F(t, N) = r - (r/K) N(t) - \int_{-\infty}^{t} k(t - s) F(N(s))\, ds$$

Where $k$ is a piecewise continuous nonnegative valued function such that $\int_0^\infty k(s)\,ds = 1$ with some appropriate function $F$, we will get Volterra's model of a population which pollutes its environment and pollution has an accumulative and toxic effect; such an equation is of the form

$$\frac{dN(t)}{dt} = N(t) = \left\{ r - (r/K) N(t) - \int_{-\infty}^{t} k(t-s) F(N(s))\,ds \right\} \quad (1.4)$$

Which has been studied in different formats by several authors (Cushing [12], Cohen *et al.* [10], and Gopalasamy [16])

One of the significant advantages of deriving various models from the prototype (1.2) is that if the initial; conditions (initial population sizes) are nonnegative then the nonnegativity of the population size $N(t)$ for $t > 0$ follows from the fact that solutions of (1.2) satisfy

$$N(t) = N(0) \exp\left[\int_0^t f(s, N(s))ds\right]$$

Which implies that $N(t) \geq 0$ if $N(0) \geq 0$.

While the time delayed logistic equation (1.3) explains some of the observed oscillations of certain populations, its derivation has not been entirely satisfactory. For instance it has been pointed out by Ricklefs ([27], page 488) that "time lags might be expected to occur primarily in stabilizing processes involving production rather than death because death is an immediate response to environmental change". Recently Cushing [13] has examined the effects of response and maturation delays in a number of single and multi system models; as it has been noted by Cushing [13], the models involving delayed reproduction or not necessarily derivable from the Kolmogorov type systems of the form (1.2) by suitably defining the average growth rate $(1/N)(dN/dt) = f(t, N(t))$. While Cushing [13] has been mostly concerned with the possibility of delay induced oscillations we will consider a class of delays in reproduction and interspecific interaction which maintain the stability of the systems.

Our starting point to drive models with delayed reproduction is the following balance equation which assumes that there is no immigration or emigration. Account, particle, atom, molecule, proposition, problem et al

$$\frac{dN(t)}{dt} = \text{birth rate} - \text{death rate}. \quad (1.5)$$

Now for instance if we consider a population of adult flies then the birth rate at time $t$ depends on the adult population at time $t^{-}\ \tau$ where $\tau$ is the time required for the larvae to become adult flies. If the birth and death rates are governed by density dependent factors then we have from (1.5) the following time delayed model

$$\frac{dN(t)}{dt} = b\big(N(t^{-}\tau)\big) - m(N(t))$$

(1.6)

Where the functions $b$ and $m$ denote the density dependent birth and death rates of Account, particle, atom, molecule, proposition, problem et al respectively. If the time delay is continuously distributed then one can consider in the place of (1.6) an equation of the form

$$\frac{dN(t)}{dt} = b\int_{-\infty}^{t} k\{(t-s)N(s)ds - m(N(t))$$

Or equivalently

$$\left\{\frac{dN(t)}{dt} = b\int_{0}^{\infty} k(s)\right\}N(t-s)ds - m(N(t))$$

(1.7)

Where the delay kernel $k$ denotes the weight of the delayed of the delayed effects.

It is not always obvious that suitable nonnegative initial conditions for (1.6) and (1.7) will imply the nonnegativity of the solutions of (1.6) and (1.7). The purpose of this article is to show that in two species systems of competitive and prey- predator communities, delayed reproduction and delayed interactions donot destabilise the systems if the self- regulating negative feedback effects are sufficiently strong compared to interspecific interactive effects and if the self regulating effects are realised with no time delays. When we say that time delays do not destabilise, we will mean that delays do not render an asymptotically (locally) stable steady state unstable in such a manner that the loss of stability either leads

To a delay – induced bifurcation to persistent and undamped oscillations (as in the case of Hopf- bifurcation) or the relevant linear variational system has unbounded solutions; that is a mathematical instability is not induced by delays in the sense, that the roots of the characteristic equation are pushed to the imaginary axis or to the right of it on the complex plane. However delays can make an otherwise stable system less stable which means that the rate of convergence to zero of solutions of the variational system can decrease with an increase of delay and the variational system can even oscillate (perhaps violently) before converging to the trivial state. Such a reduction in stability can

happen even when all the roots of the characteristic equation have negative real parts. For our purposes such a reduced stability is still stability. In order to see the effects of delays on such less stable systems we will propose below a method of estimating the rate of decay of local perturbations based on the variational system and indirectly relying on the characteristic equation. For a special class of two species models with a single delay we have obtained (Gopalsamy [16]) estimates of decay rates (or return times) usingthe characteristic equation; such a procedure is almost impossible in the case of many species systems with several delays. The method we propose in Section 6 below for the estimate of decay rates can be adapted to *n*- species systems with $n^2$ discrete delays.

Let $x_1(t)$ and $x_2(t)$ denote the population sizes (or biomasses) of two systems competing for a common pool of resources.. Let $b_i$ and $m_i$ $(i = 1, 2)$ denote the respective density dependent birth (potential) and death rates so that in the absence of any time delays the population sizes are governed by (of Account, particle, atom, molecule, proposition, problem et al). In addition as usual, we assume that the system studied follow Haimovici's equations, for which the dove tailing explanation has been appended

$$\left.\begin{aligned}\frac{dx_1(t)}{dt} &= b_1\left(x_1(t)\right) - m_1(x_1(t)), x_2(t))\\ \frac{dx_2(t)}{dt} &= b_2\left(x_2(t)\right) - m_2(x_1(t)), x_2(t))\end{aligned}\right\} \quad (2.1)$$

In order to make the system (2.1) represent a competition system we will make the following assumptions on the birth and death rates; $b_1$, $m_1$ $(i = 1, 2)$ are continuous with continuous partial derivatives for all $x_1 \geq 0$ $(I = 1, 2)$; also we assume (see for instance Rescigno and Richardson [26]) $\frac{\partial b_i}{\partial x_1} > 0, \frac{\partial m_1}{\partial x_j} > 0$ for $x_i > 0$ $(i, j = 1, 2)$, (2.2)

(x) $b_1(0) = 0$, $m_1(0, x_2) \equiv 0, b_2(0) = 0$, $m_2(x_1, 0) \equiv 0$, (2.3)For some $x_{1,}^* > 0, x_{2,}^* > 0$ we have

$$\left.\begin{aligned}b_1(x_1^*, -m_1(x_1^*, 0) &= 0,\\ b_2(x_2^*, -m_1(0, x_2^*) &= 0,\end{aligned}\right\} \quad (2.4)$$

(xi) There exist positive constants $\delta_1$, $\delta_2$ such that

$b_1(\delta_1) - m_1(\delta_1, x_2) < 0$,

$b_2(\delta_2) - m_2(x_1, \delta_2) < 0$, for $x_1 \geq 0$, $x_2 \geq 0$,

(2.5)

(xii) There exists a point $(\alpha, \beta), \alpha > 0, \beta > 0$ such that

$$b_1(\alpha) - m_1(\alpha, \beta) = 0 = b_2(\beta) - m_2(\alpha, \beta).$$

(2.6)

The conditions on $b_i$ in (2.2) mean that the potential (not average) birth rates are density dependent and any negative density dependent feedbacks to the opened or birth rate of Account, particle, atom, molecule, proposition, problem et al rates are included in the destruction rates; the conditions on $m_i$ indicate intraspecific and interspecific confederation needs and necessities. (2.3) implies that (0, 0) is a (trivial) steady state of the system (2.1) as it is in the case of all ecological models; (2.4) means that in the absence of either of the two species the other one has a nontrivial steady state; (2.5) implies that each of the two populations cannot grow unboundedly since when $x_1(t) > \delta_i, \frac{dx_1}{dt} < 0$. (2.6) means that there exists a nontrivial steady state $(\alpha, \beta)$ in the interior of the positive quadrant of the $(x_1, x_2)$- plane. These assumptions are standard and are minimal in the opinion of the author.

An immediate example of (2.1) –(2.6) is the familiar Volterra –Lotka competition system described by

$$\frac{dx_1(t)}{dt} = r_1 x_1(t) - a_{11} x_1^2(t) - a_{12} x_1(t) x_2(t),$$

$$\frac{dx_2(t)}{dt} = r_2 x_2(t) - a_{21} x_1(t) x_2(t) - a_{22} x_2^2(t),$$

(2.7)

Where $r, a_{ij}, (i, j = 1,2)$ are positive constants satisfying the condition

$$\frac{a_{11}}{a_{21}} > \frac{r_1}{r_2} > \frac{a_{12}}{a_{22}}.$$

One of the first questions to be resolved for (2.1) is the following: if $x_i(0) > 0$, $x_2(0) > 0$ does it follow that the corresponding solutions of (2.1) satisfy the condition $x_1(t) \geq 0$, $x_2(t) \geq 0$ for those $t$ for which a solution of (2.1) exists? In the $(x_1, x_2)$ phase plane of (2.1) we have (0, 0), $(x_1^*, 0)$, $(0, x_2^*)$, $(\alpha, \beta)$ as the steady states of (2.1) and also it will follow from (2.3) that the co ordinate axes

$x_1 = 0$ and $x_2 = 0$ are invariant sets f (2.1) ; such invariance of the coordinate axes will imply that if $x_1(0) > 0$, $x_2(0) > 0$ then the corresponding population trajectory $(x_1(t), x_2(t))$ of (2.1) cannot reach the outside of the nonnegative quadrant $x_1 \geq 0$, $x_2 \geq 0$.

Our plan is to derive first a set of sufficient conditions for the coexistence of the two competing species and then in the next section show that under the same conditions the two species can coexist if there are time delays in the reproduction and interspecific interactions. We recall that the two competing species described by (2.1) are said to coexist if and only if the nontrivial steady state $(\alpha, \beta)$ is (at least locally) asymptotically stable. To examine such a stability of $(\alpha, \beta)$ we derive the corresponding linear variational system for the perturbations $X_1, X_2$ where

$$x_1(t) \equiv \alpha + X_1(t), \qquad x_2(t) \equiv \beta + X_2(t)$$

Such a variational system is found to be

$$\left.\begin{aligned} \frac{dX_1(t)}{dt} &= \beta_{11} X_1(t) + \beta_{12} X_2(t), \\ \frac{dX_2(t)}{dt} &= \beta_{21} X_1(t) + \beta_{22} X_2(t), \end{aligned}\right\} \qquad (2.8)$$

Where

$$\beta_{1J} = \left(\frac{\partial b_1}{\partial x_J} - \frac{\partial m_1}{\partial x_j}\right) \text{ evaluated at } (\alpha, \beta) \qquad (i, j = 1,2).$$

(2.9)

The steady state $(\alpha, \beta)$ of (2.1) is locally asymptotically stable if and only if the trivial solution (0, 0) of (2.8) is asymptotically stable. We can derive the following result from elementary considerations.

In the system (2.1)-(2.6) assume the following:

$$\left.\begin{aligned} \frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1}, \\ \frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1}, \end{aligned}\right\} \qquad \text{at } (\alpha, \beta)$$

(2.10)

Then the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (2.1) – (2.6).

The proof follows easily if we note that the characteristic equation associated with the linear system (2.8) is given by

Det $\begin{pmatrix} \lambda-\beta_{11} & -\beta_{12} \\ -\beta_{21} & \lambda-\beta_{22} \end{pmatrix} = 0$

Or equivalently,

$$\lambda^2 - \lambda\,(\beta_{11} + \beta_{22}) + \beta_{11}\beta_{22} - \beta_{12}\beta_{21} = 0. \qquad (2.11)$$

It follows from (2.2) and (2.10) that,

$$\beta_{11} < 0, \beta_{11}\beta_{22} - \beta_{12}\beta_{21} > 0$$

And hence the roots of (2.11) have negative real parts implying the asymptotic stability of the trivial solution of (2.8)

If apply the condition (2.10) to the Volterra-Lotka model (2.7) we note that when a nontrivial steady state say $(\alpha, \beta)$ exists for (2.7), (2.10) leads to

$$a_{11}\alpha > a_{21}\,\beta \; and \; a_{22}\,\beta > a_{12}\,\alpha,$$

Which together imply $a_{11}\,a_{22} > a_{12}a_{21}$. If $(\alpha, \beta)$ exists it is known that $a_{11}\,a_{22} > a_{12}a_{21} > 0$ is a sufficient condition for the (local) asymptotic stability of $(\alpha, \beta)$.

The condition (2.10) which guarantees the coexistence of competing populations in (2.1)-(2.6) simply means that the intraspecific self – regulating feedback effects $\partial m_i / \partial x_1\ (i = 1, 2)$ bare higher than its own positive feedback $\partial b_1 / \partial x_1\ (i = 1, 2)$ together with its negative (competitive) effects on its competitor $\partial m_I / \partial x_J\ (i,\ j = 1, 2; i \neq j)$.

## 8. Coexistence under delayed production and interactions, transactions

We will now establish that two competing species with a sufficiently strong intraspecific regulation can coexist in the presence of delayed reproduction and delayed interspecific interactions. It is well known that if there are delays in the self – regulating intraspecific effects, stationary coexistence may not always be possible (Gopalsamy [16], Gopalsamy and aggarwala [15], Shibata and Saito [28]). Delayed reproduction is more general and widespread in most biological populations. To build a competition model with such delays we proceed as follows.

Let $\tau_{ij}\ (i, j = 1, 2)$ be a set of nonnegative constants with $\tau = \max \tau_{ij}\ , i, j = 1, 2,$ and suppose that the two ompeting species display delayed reproduction and delayed interspecific interactions with no delays in the self – regulating negative feedback effects. Such a competition system in a constant environment can be modeled by the autonomous delay differential system

$$\frac{dx_1(t)}{dt} = b_1\big(x_1(t-\tau_{11})\big) - m_1\big(x_1(t), x_2(t-\tau_{12})\big),$$

$$\frac{dx_2(t)}{dt} = b_2\big(x_2(t-\tau_{22})\big) - m\big(x_1(t-\tau_{21}), x_2(t)\big), \qquad (3.1)$$

Where both the creation and destruction rates and the morality rates $b_1$ and $m_1\ (i = 1, 2)$ satisfy the same conditions (2.2)-(2.6) as in the previous section. Along with (3.1) we will

Suppose that the initial population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) sizes are provided in the form

$$x_i(s) = \phi_i(s) \geq 0, \qquad s \in [-\tau, 0], \qquad i = 1,2. \tag{3.2}$$

(Note that the specification of $\phi_i$ need not be on the entire interval$[-\tau, 0]$; but such a hypothesis is convenient.) Since (3.1) is not of the standard Kolmogorov-type, we have to establish that the solutions of (3.1) will remain nonnegative for nonnegative initial conditions.

To prove the nonnegativity of solutions of (3.1)-(3.2), let $\left\{\psi_1^{(n)}(s)\right\}, s \in [-\tau, 0], i = 1,2; (n = 1,2, \ldots)$ be a sequence of strictly positive continuous functions such that

$$\lim_{n\to\infty} \psi_1^{(n)}(s), = \psi_1(s), \qquad s \in [-\tau, 0], \qquad i = 1,2 \tag{3.3}$$

Let $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ be the solution of (3.1) corresponding to the initial condition

$$\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \left\{\psi_1^{(n)}(s), \psi_2^{(n)}(s)\right\}, \quad s \in [-\tau, 0],. \tag{3.4}$$

Consider the solution $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ for $t \in [0, \tau]$ (assuming it exists on the whole of $[0, \tau]$). Suppose $x_1^{(n)}$ does not remain nonnegative for all $[0, \tau^*]$; then there exists a $t^* = [0, \tau^*](\tau^* =$ the positive minimum of $\tau_{ij}$ $(i, j = 1, 2))$ for which

$$x_1^{(n)}(t^*) = 0 \ and \ x_1^{(n)}(t)(t^* - \tau_{11})) > 0.$$

It will follow from (3.1) and the positivity of the initial condition that

$$\frac{dx_1^{(n)}(t^*)}{dt} = b_2\left(x_1^{(n)}(t^* - \tau_{11})\right) > 0 \tag{3.5}$$

Where we have used the properties of $b_i$ and $m_i$ from (2.2)-(2.6); it is found that (3.5) contradicts the definition of $t^*$. Thus $x_1^{(n)}(t) > 0$ on $[0, \tau^*]$; similarly $x_2^{(n)}(t) > 0\ on\ [0, \tau^*]$. We can repeat this procedure for intervals of the form $[\tau^*, 2\tau^*], [2\tau^*, 3\tau^*]\ etc$. Thus it will follow that so long as $(x_1^{(n)}(t), x_2^{(n)}(t))$ exists, we have

$$x_1^{(n)}(t) \geq 0\ x_2^{(n)}(t) \geq 0.$$

If we now consider the limits as $n \to \infty$, we get

$$\lim_{n\to\infty}\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \{x_1(t), x_2(t)\} \tag{3.2}$$

With $x_1(t) \geq 0, , x_2(t) \geq 0$ where $(x_1(t), x_2(t))$ is the solution of (3.1)-(3.2) and this is a consequence of the continuous dependence of solutions of (3.1) on initial conditions (Hale [17], page 41).

A second question for (3.1) with (2.2)-(2.6) is whether solutions of (3.1) exist for all $t > 0$? suppose now that solutions of (3.1) do not exist for all $t$ ; then there exists a $t_1 > 0$ such that at least for one of either of $x_1$ *or* $x_2$ we have

$$\lim_{t \to t_1 -} x_i\ (t) = \infty. \tag{3.7}$$

To be specific let us suppose that $\lim_{t \to t_1 -} x_i\ (t) = \infty$; let $t_2$ be the first instant of time for which $x_1(t_2) = \delta_1, t_2 < t_1$. It will follow from (3.1) that

$$\frac{dx_1(t_2)}{dt} = b_1\big(x_1\ (t_2 - \tau_{11})\big) - m_1\big(x_1(t_2), x_2\ (t_2 - \tau_{12})\big)$$

$$< b_1\big(x_1\ (t_2)\big) - m_1\big(x_1(t_2), x_2\ (t_2 - \tau_{12})\big)$$

$$< b_1(\delta_1) - m_1\big(\delta_1, x\ (t_2 - \tau_{12})\big) < 0 \tag{3.8}$$

Which contradicts the definition of $t_2$. Thus solutions of (3.1), (2.2) – (2.6) exist for all $t$ .

We are now ready to consider the effects of time delays in reproduction and interspecific interactions on the coexistence of the two competing species. The steady state $(\alpha, \beta)$ of (2.1) is again a steady state of (3.1). To examine the (local) asymptotic stability of $(\alpha, \beta)$ in (3.1) we let

$$x_1(t) = \alpha + Y_1(t), \quad x_2(t) = \beta + Y_2(t)$$

In (3.1) and derive that the linear variational system governing the perturbations $Y_1, Y_2$ is given by

$$\frac{dY_1\ (t)}{dt} = -m_{11}Y_1(t) + b_{11}Y_1\ (t - \tau_{11}) - m_{12}Y_2\ (t - \tau_{12}),$$

$$\frac{dY_2\ (t)}{dt} = -m_{22}Y_2(t) + b_{21}Y_1\ (t - \tau_{21}) - m_{22}Y_2\ (t - \tau_{22}),$$

Where

$$m_{ij} = \frac{\partial m_1}{\partial x_j}, \quad b_{ii} = \frac{\partial b_i}{\partial x_i}\ at\ (\alpha, \beta). \tag{3.10}$$

It is well known (Bellman and Cooke [3], page 336) that the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) if the trivial solution of the variational system (3.9) is asymptotically stable. We formulate our first result in the following form whose proof is given in the appendix.

*THEOREM 3.1. Assume that in the two systemic types competition model system (3.1) with delayed density dependent birth and death rates $b_i$ and $m_i$ $(i = 1, 2)$ satisfy the conditions (2.2)-(2.6). Suppose that the conditions (2.10) hold. Then the nontrivial steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) whatever the delays be $\tau_{ij}$ $(i = 1, 2)$.*

Let $y_1(t) and\ y_2(t)$ respectively denote the population sizes (or biomasses) of a prey species and a predator species which feeds exclusively on the single prey species in a temporally constant environment. In the format of (1.5) the following system will describe the rates of change of the two population sizes:

$$\frac{dy_1(t)}{dt} = f_1(y_1(t)) - g_1(y_1(t), y_2(t)),$$

$$\frac{dy_2(t)}{dt} = f_2(y_1(t)) - y_2(t) - g_2(y_2(t)), \quad (4.1)$$

Where $f_i\ and\ g_i\ i = 1,2)$ respectively denote the birth and death rates of the two species; in order to make (4.1) denote the dynamics of a prey – predator species we will assume that $f_i\ and\ g_i$ satisfy the following conditions (see for instance Rescigno and Richardson [26]); $f_i\ and\ g_i$ are continuous with continuous partial derivatives for all $y_i \geq 0\ (i = 1, 2); also\ we\ require$

$$\frac{\partial f_1}{\partial y_1} > 0; \quad \frac{\partial g_1}{\partial y_1} \geq 0; \quad \frac{\partial g_1}{\partial y_2} \geq 0,$$

$$\frac{\partial f_2}{\partial y_1} \geq 0; \quad \frac{\partial g_2}{\partial y_2} \geq 0; \quad \frac{\partial g_2}{\partial y_2} \geq 0, \quad \text{Where } y_1 \geq 0, y_1 \geq 0, \quad (4.2)$$

(xiii)

$$f_1(0) = 0, \quad g_1(0, y_2 \equiv 0,$$

$$f_2(0, y_2) = 0 = f_2(y_1, 0); \quad g_2(0) = 0, \quad (4.3)$$

(xiv) For some constant $y_1^* > 0$ we have

$$f_1(y_1^*) - g_1(y_1^* 0) = 0, \quad (4.4)$$

(xv) There exists a point $(\alpha^*\beta^*), \alpha^* > 0, \beta^* > 0$ such that

$$f_1(\alpha^*) - g_1(\alpha^*, \beta^*) = 0,$$

$$f_2(\alpha^*, \beta^*) - g_2(\beta^*) = 0, \quad (4.5)$$

(xvi) There exist positive constants $n_1, n_2$ such that

$$f_1(\eta_1) - g_1(\eta_1, y_2) < 0, y_2 \geq 0,$$

$$f_2(\eta_1, \eta_2) - g_2(\eta_2) < 0 \quad (4.6)$$

The conditions (4.2) mean that the system (4.1) is of the oppugnacy systems type; (4.3) will imply that (0,0) is a trivial steady state of (4.1); (4.4) means that in the absence of the predator, the prey species has a nontrivial (positive) steady state;

(4.5) implies that (4.1) has a nontrivial steady state in the interior of the positive quadrant of the state species of (4.1); (4.6) can be shown (see below) to imply that no credit type and debit type system types can grow unboundedly.

An example of (4.1)-(4.6) is the familiar Volterra-Lotka prey-predator system

$$\frac{dy_l(t)}{dt} = r_l y_l(t) - [a_{ll} y_l^2(t) + a_{12} y_l(t) y_2(t)],$$

$$\frac{dy_2(t)}{dt} = a_{12} y_1(t) y_2(t) - [r_2 y_2(t) + a_{22} y_2^2(t)],$$

Where $r_{l,}\ a_{ij}\ (i,j=1,2)$are positive costants such that

$$a_{21}r_1 - r_2a_{11} > 0.$$

As in the previous case of competition we have to verify that when $y_1(0) > 0, y_2(0) > 0$ the corresponding solution of (4.1) will be nonnegative for $t \geq 0$ for which the solution exist. Since the co-ordinate axes $y_1 = 0, y_2 = 0$ are invariant sets for (401) by virtue of (4.3), it will follow that the solution of (4.1) starting from the interior of the positive quadrant cannot reach the outside of that quadrant. This is a consequence of the invariance of the co-ordinate axes for (4.1)

To prove that solutions of (4.1)-(4.6) exist for all t $\geq 0$ let us suppose that $y_l 0\ and\ y_2(0) satisfy\ 0 < y_l\ (0)$ satisfy $0< y_l(0) < \eta_l, 0 < y_2(0) < \eta_2$ ( on the other hand if either one of $y_l(0) \geq \eta_l and\ y_2(0) \geq \eta_l$ or both hold, it will follow that the corresponding $y_l(t)$ will decrese by virtue of (4.6) at least for small $t \geq 0,$ ). Suppose $y_l(t)$ does not exist for all $t \geq 0$; then for some $t_l < \infty\ lim_{t\to t_l-} y_l(t) = \infty$and let $t_2 > 0$ be the first instant for which $y_1(t_2) = \eta_1$; then we have from (4.1) and (4.6) that

$$\frac{dy_1(t_2)}{dt} = f_1(\eta_1) - g_1(\eta_{1,}\ y_2) < 0\ (assumimg\ y_1(0) < \eta_1)$$

Which contradicts the definition of $t_2$ showing that $y_1(t)$ exists for all $t \geq 0$; if $y_2(t)$ is not defined for all $t \geq 0$, then for some $t_3 \geq \infty$ we will have $lim_{t-t_3-} y_2(t) = \infty$;now let $t_4$ be the first instant for which $y_2(t_4) = \eta_2$; then from the properties of $f_2$ we have

$$\frac{dy_2(t_4)}{dt} = f_2(y_1(t_4), \eta_2) - g_2(\eta_2) \leq f_2(\eta_1, \eta_2) - g_2(\eta_2)$$

$$< 0 (by\ (4.6) assuming\ 0 < y_1(0) < \eta_1, 0 < y_2(0) < \eta_2)$$

Which again contradicts the definition of $t_4$. Thus $y_2(t)$ is defined for all $t \geq 0$.

We will now derive a set of sufficient condition for the (local) asymptotic stability of the steady state $(\alpha^*, \beta^*)$ of (4.1). in the next section we will show that under the same conditions and in the presence of various delays , the steady state $(\alpha^*, \beta^*)$remains (locally) asymptotically stable.

To examine the asymptotic stability of the nontrivial steady state $(\alpha^*, \beta^*)$ in (4.1) we let

$$y_1(t) = \alpha^* + U_1(t), \qquad y_2(t) = \beta^* + U_2(t)$$

In (4.1) and derive the liner variational system in the perturbations $U_1, U_2$ as follows:

$$\left.\begin{aligned} \frac{dU_1(t)}{dt} &= \frac{\partial f_1}{\partial y_1}U_1(t) - \frac{\partial g_1}{\partial y_1}U_1(t) - \frac{\partial g_1}{\partial y_2}U_2(t), \\ \frac{dU_2(t)}{dt} &= \frac{\partial f_2}{\partial y_1}U_1(t) - \frac{\partial f_2}{\partial y_2}U_2(t) - \frac{\partial g_2}{\partial y_2}U_2(t), \end{aligned}\right\} \quad (4.7)$$

Where all the partial derivatives in (4.7) are evaluated at $(\alpha^*, \beta^*)$. The steady state $(\alpha^*, \beta^*)$ of (4.1) is (locally) asymptotically stable if and only if the trivial solution of (4.7) is asymptotically stable which will be the case if and only if the roots of the characteristic equation

$$\text{Det} \begin{pmatrix} \lambda - (f_{11} - g_{11}) & g_{12} \\ -f_{21} & \lambda - (f_{22} - g_{22}) \end{pmatrix} = 0,$$

Where $f_{ij} = \frac{\partial f_1}{\partial y_l}$, $g_{ij} = \frac{\partial g_1}{\partial y_j}$, at $(\alpha^*, \beta^*)$ (4.8)

$(i,j = 1, 2)$

Have negative real parts. If $\lambda_1, \lambda_2$ are the riits of (4.8) then

$$\lambda_1 + \lambda_2 = (f_{11} - g_{11}) + (f_{22} - g_{22}), \quad (4.9)$$

$$\lambda_1 \lambda_2 = (f_{11} - g_{11})(f_{22} - g_{22}) + f_{21} g_{12}.$$

The proof of the following result is an immediate consequence of (4.7)-(4.9)

THEOREM 4.1 *In the CREDIT-DEBIT system (4.1)-(4.6) assumes that the partial derivatives satisfy the following additional conditions:*

$$\frac{\partial g_1}{\partial y_1} > \frac{\partial f_1}{\partial y_1} + \frac{\partial f_2}{\partial y_1},$$

$$\frac{\partial g_2}{\partial y_2} > \frac{\partial f_2}{\partial y_2} + \frac{\partial g_1}{\partial y_2}, \quad \text{Evaluated at } (\alpha^*, \beta^*), \quad (4.10)$$

*Then the steady state $(\alpha^*, \beta^*)$, is (locally) asymptotically stable for the prey- predator system.*

## 8. Delays in reproduction, predation and gestation

A number of authors have proposed and analysed several models of prey- predator systems with time delays. Most of the models have been derived from the so called models of the Kolmogorov-type f (1.2) with time delays incorporated in the average growth rates of the prey and predator species. Volterra [32] proposed the following model for a prey-predator system. The equations are used for the formulation of Credit-debit system of atoms, accounts, economic variable, cosmological variable, astrophysical; variables, and qubit systems for which the model is being applied.

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \delta_1 N_1(t) - \int_{\infty}^{t} F_1(t-s) N_2(s) ds],$$

(5.1)

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 + \delta_2 N_1(t) - \int_{\infty}^{t} F_2(t-s)N_1(s)ds],$$

Where $r_1, r_2, \delta_1, \delta_2$ *are positive constant and* $F_1, F_2$ is nonnegative continuous delay kernel defined and integrable on $[0, \infty]$. (5.1) has no predator self-regulating negative feedback. Brelot [6] considered a modified format of (5.1) in the form

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \lambda_1 N_1(t) - \delta_1 N_2(t) - \int_{-\infty}^{t} F_1(t-s)N_2(s)ds],$$

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 - \lambda_2 N_2(t) - \delta_2 N_1(t) - \int_{-\infty}^{t} F_2(t-s)N_1(s)ds], \quad (5.2)$$

Which has term representing intraspecific competition among both predator and prey species. Under suitable condition on the various parameters in (5.2) one can show that (5.2) has a (locally) asymptotically stable steady state. A number of intergrodifferential equation models of the types (5.1) and (5.2) have been investigated by Cushing [12]

Starting from Hutchinson's (1948)[19] delay logistic equation, May[23] has proposed the following system

$$\frac{dN_1(t)}{dt} = rN_1(t)\left[1 - \frac{N_1(t-\tau)}{K}\right] - \alpha N_1(t)N_2(t),$$

$$\frac{dN_2(t)}{dt} = -bN_2(t) + \beta N_2(t)N_1(t), \quad (5.3)$$

Where r, $\tau$, K, $\alpha, \beta, b$ are positive constants; (5.3) contains a single discrete delay; one can modify (5.3) and incorporate a continuously distributed delay in it so that (5.3) can be written as (May [23])

$$\frac{dN_1(t)}{dt} = rN_1(t)[1 - \int_{-\infty}^{t} F(t-s)N_1(s)ds] - \alpha N_1(t)N_2(t),$$

$$\frac{dN_2(t)}{dt} = -bN_2(t) + \beta N_1(t)N_2(t). \quad (5.4)$$

The first model of a prey-predator type which departs from the Kolmogorov-type formulation is due to Wangersky and Cunningham [33] who considered the model system

$$\frac{dN_1(t)}{dt} = a_1N_1(t) - b_1N_1^2(t) - c_1N_1(t) - N_2(t),$$

$$\frac{dN_2(t)}{dt} = a_2N_2(t) + c_2N_1(t-\tau)N_2(t-\tau), \quad (5.5)$$

Where $a_1, a_2, b_1, c_1, c_2, \tau$ are positive constants; (5.5) means that a duration of $\tau$ time units elapses when an individual prey is killed and the moment when the corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) Recently Cushing [13] has considered a model of the form

$$\frac{dN_1(t)}{dt} = rN_1(t)\left[1 - \frac{N_1(t)}{k} - \alpha N_2(t)\right],$$

$$\frac{dN_2(t)}{dt} = -d_2^*(t) + bN_1(t)\int_0^\infty \delta(a)N_2(t-a)e^{-d*a}\,da,$$

(5.6)

Where $r, k, \alpha, d^*, b$ are positive constants and $\delta$ is related to the age dependent fecundity of the predator species.

With the models (5.1)-(5.6) in the background let us now consider a time delayed model similar to (4.1) in the following form:

$$\frac{dN_1(t)}{dt} = f_1\big(N_1(t-\tau_{11})\big) - g_1\big(N_1(t), N_2(t-\tau_{12})\big),$$

$$\frac{dN_2(t)}{dt} = f_2\big(N_1(t-\tau_{21})\big), N_2(t-\tau_{12}) - g_1(N_2(t)),$$

(5.7)

Where $\tau_{ij}$ $(i,j=1,2)$ are nonnegative constants with $\tau = \max \tau_{ij}, i,j = 1,2$; the birth and death rate functions $f_1, f_2, g_1, g_2$ satisfy the same conditions as in (4.1)-(4.6). Along with (5.7) we specify the initial conditions as follows:

$$N_1(s) = \phi_1(s) > 0,$$

$$s \in [-\tau, 0].$$

$$N_1(s) = \phi_1(s) > 0,$$

We3 will assume that in addition to (4.2) we have for (5.7) that

$$\frac{\partial f_2}{\partial N_1} > 0, \frac{\partial f_2}{\partial N_2} > 0, f_2(N_2, N_2) > 0 \; whenever \; N_1 > 0, \; N_2 > 0. \; (5.9)$$

It is not difficult to show by the methods previously used for the competition system that whenever $N_1(s) > 0$ and $N_2(s) > 0$ for $s \in [-\tau, 0]$, solution of (5.7) exist for all $t \geq 0$ and remain nonnegative for all $t \geq 0$. also $(\alpha^*, \beta^*)$ of (4.5) is a steady state of (5.7).

We will discuss the form and meaning of the delays in (5.7) in the next section. The effects of the various delays in (5.7) on the (local) asymptotic stability of $(\alpha^*, \beta^*)$ can be derived from the following theorem whose proof is quite similar to that of Theorem 3.1 and hence the details of the proof are omitted.

THEOREM 5.1 *In the time delayed prey-predator system(* corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) *) (5.7) let the conditions (4.2)-(4.6), (5.8) and (5.9) hold for the birth and death rates; furthermore suppose the condition (4.10) hold. Then for all nonnegative delays in (5.7), its nontrivial steady state $(\alpha^*, \beta^*)$ is (locally) asymptotically stable.*

## 6. Estimation of decay rates

We will first derive a convention sufficient condition for the asymptotic stability of the trivial solution of a linear system of the form

$$\left.\begin{aligned} \frac{dY_1(t)}{dt} &= -m_{11} Y_1(t) + b_{11} Y_1(t-\tau_{11}) - m_{12} Y_2(t-\tau_{12}), \\ \frac{dY_2(t)}{dt} &= -m_{22} Y_1(t) + m_{21} Y_1(t-\tau_{21}) - b_{22} Y_2(t-\tau_{22}), \end{aligned}\right\} \; t > 0, \quad (6.1)$$

Where $m_{1J}, b_{i1} \; (i, j = 1,2)$ are positive constants while $\tau_{1J} (i, j = 1,2)$ are non negative with $\tau = \max(\tau_{1J})(i, j = 1,2$.as is shown in the appendix, the following system dominates (6.1*)*

$$\left.\begin{aligned} \frac{dZ_1(t)}{dt} &= -m_{11} Z_1(t) + b_{11} Z_1(t-\tau_{11}) - m_{12} Z_2(t-\tau_{12}), \\ \frac{dZ_2(t)}{dt} &= -m_{22} Z_1(t) + m_{21} Z_1(t-\tau_{21}) - b_{22} Z_2(t-\tau_{22}), \end{aligned}\right\} \; t > 0, \quad (6.2)$$

$Z_i(s) = |Y_i(s)|, \qquad i = 1,2 \;\; for \; s \leq 0 \;\; for \; which \; Y_i(s) \; is \; defined.$

Thus the asymptotic stability of the trivial solution of (6.2) implies that of (6.1). it is also shown in the appendix that the solutions of (6.2) are nonnegative for all t $\geq 0$. we will use the following vector-matrix notation:

$$Z = (Z_1, Z_2)^{\tau}, \qquad ||Z|| = |Z_1(t)| + |Z_2(t)| = Z_1(t) + Z_2(t)$$

$$B = \begin{bmatrix} b_{11} & m_{12} \\ m_{21} & b_{12} \end{bmatrix}, \quad ||B|| = max \begin{matrix} b_{11} + & m_{21} \\ b_{21} + & m_{12} \end{matrix}, \tag{6.3}$$

$$A = \begin{bmatrix} -m_{11} & 0 \\ 0 & -m_{22} \end{bmatrix}.$$

The logarithmic norm or the measure of any square matrix B, denoted by $\mu(B)$is defined as follows:

$$\mu(B) = \lim_{\theta \to 0+} \frac{/\!/I + \theta B/\!/ - 1}{\theta}$$

($I$ being an identity matrix of the same size as $B$).The matrix measure compatible with the norms in (6.3) is given by

$\mu(B) = \max_J [b_{jj} + \sum_{\substack{i=1 \\ i \neq J}}^{2} | b_{jJ} | ]$ $B=(b_{1J})$.

We refer to Coppel [11]or Vidyasager [31]for a number of properties of matrix measure We have shown in the appendix that the derivative of $H(Z_1,Z_2)(t)$(seeA.10)satisfies the relation

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq [-m_{11}+b_{11}+m_{21}]Z_1(t)+[-m_{22}+b_{22}+m_{12}]Z_2(t)$, and hence we have

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq \{\mu(A) + /\!/B/\!/\}\{ Z_1(t)+ Z_2(t)\}$

If we assume for the matrix norm $\|B\|$ as in (6.3)that

$$\mu(A) + /\!/B/\!/ < 0 \; or \; -\mu(A) > /\!/B/\!/$$

Then from the results of the appendix (Lemmas 5 and 6) it will follow that $Z_1(t)+Z_2(t) \to 0$ as $t \to \infty$ and thus we have the following:

Lemma 6.1 Assume that the Coefficients of the linear system (6.1)satisfy the condition (6.7). Then the trivial solution of(6.1)is asymptotically stable.

When (6.7)holds an estimate for the decay rate of solutions of (6.1)can be obtained as follows A necessary and sufficient condition for the asymptotic stability of the trivial solution is that all the roots of the characteristic equation

Det $G(\lambda,T_{iJ})=0$

Have negative real parts where

$G(\lambda,T_{1J})=[\begin{matrix} \frac{\lambda+m_{11}}{0} & \frac{0}{\lambda+22} \end{matrix}] - \begin{matrix} b_{11}{}^{e-\lambda T}{}_{11} & -m_{12}{}^{e-\lambda T}{}_{12} \\ -m_{21}{}^{e-\lambda T}{}_{21} & b_{22}{}^{e-\lambda T}{}_{22} \end{matrix}$

$=\lambda I - A - B(\lambda, T_{ij})$(say),

A positive number $\mathfrak{a}$(if it exists)is said to be a decay rate of solution of (6.1) if and only if every root of (6.8)satisfies the condition

$\mathrm{Re}(\lambda)+\mathfrak{a} < 0.$

When (6.11)holds the solution of (6.1)will be of the order of $e^{-at}$ and hence the supremum of all such $\mathfrak{a} > 0$ $a$ satisfying (6.11)will be a desirable parameter to estimate However such a supremum of all such (6.11)may not be possible to determine since it has to be done directly from (6.8)-(6,10)(see also Section 3).We will obtain an estimate of $\mathfrak{a}$ (not necessarily the supremum)I (6.11)by the following procedure.

Replace $\lambda$ byz-$\mathfrak{a}$ in (6.8)and define

$G(z-\mathfrak{a}, T_{IJ}) = M(z) = (z-a)I - A - B(z-\mathfrak{a}, T_{IJ})$.

By Lemma 6.1 a Sufficient condition for all roots of M(z)=0 to have negative real parts equivalently to satisfy

$\mathrm{Re}(\lambda) < -a$

is the following :

$-\mu(aI + A) > \|B_1(a, T_{ij})\|,$

Where

$$B_1(a, T_{ij}) = \begin{pmatrix} b_{11}^{\,e-\lambda T_{11}} & -m_{12}^{\,e-\lambda T_{12}} \\ -m_{21}^{\,e-\lambda T_{21}} & b_{22}^{\,e-\lambda T_{22}} \end{pmatrix}$$

Since $\mu(\mathfrak{a} I + A) = \mathfrak{a} + \mu(A)$ we have from (6.13) that

$\mathfrak{a} + \mu(A) + \|B_I(\mathfrak{a}, T_{IJ})\| < 0$

Is a sufficient condition for Re $(\lambda) < -a$If we denote the left side of (6.15)by $\mathfrak{a} f(\mathfrak{a}, T_{Ij})$ we note that

$f(0, T_{Ij}) = \|B\| + \mu(A) < 0$ $if$ (6.7)$holds$;

Al so $\partial f / \partial \mathfrak{a} > 0, f$ is continuous and monotonic in $(\mathfrak{a})$;$\to \infty as\ \mathfrak{a} \to \infty$; thus by the continuity of $f$ on $\mathfrak{a} \geq 0$, there exists an interval of the from $[0, ^{*})$such that

$f(0, T_{Ij}) < 0$ $for\ \mathfrak{a} \in [0, \mathfrak{a}^{*})$

And hence

$f(\mathfrak{a}, T_{Ij}) = 0$

Thus an upper bound $\mathfrak{a}^*$ of $\mathfrak{a}$ for which we have $f(\mathfrak{a}, T_{Ij}) < 0$ $\mathfrak{a} \in [0, \mathfrak{a}^*)$ can be obtained as the unique (by the monotonicity of $f$ on $\mathfrak{a}$) positive root of$\max \{ b_{11}^{e-\lambda T}{}_{11} - m_{12}^{e-\lambda T}{}_{21} + \mathfrak{a}\text{-}\min \{ m_{11} \}_{=0}$

$b_{22}^{e-\lambda T}{}_{22} - m_{21}^{e-\lambda T}{}_{12} \quad m_{22} \}$

Where the maximum in (6.18) is taken over $b_{11}, b_{22}, m_{12}, m_{21}$ and $T_{ij} (i, = j\ 1,2)$. A numerical procedure can be devised to solve equations of the form (6.18) in concrete cases. If one is interested in a smaller bound on the decay rate it can be obtained as the unique root of

$\| B \| e^{ar} + a + \mu(A) = 0, \tau = \max_{i,J} T_{ij}$. (6.19)

It will follow from (6.19) than an increase in T

Will reduce the value of the root of (6.19). We remark that since our estimate $a^*$ of dacay rate is based on a sufficient condition of the form (6.7), we can conclude that the solotions of (6.1) will converge at least as fast as $e^{-a*t}$ where $a^*$ is the unique root of either of (6.18) or (6.19).

## Implications

Authors have formulated a class of two systems model systems staring from the elementary principal that the rate of change of a population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) equals the creation rate minus the destruction rate. This is a departure from the conventional model building in population ecology where one usually starts from the Malthusian formulation, now known as the Kolmogorov formulation in the nonlinear context. In our analysis of the models proposed we have obtained sufficient conditions for the global existence of nonnegative solutions of the model systems (this is trivial in the Kolmogorov formulation) and the other desirable properties such as the existence of suitable steady states. The two classes of models proposed for competition and debt contain intraspecific self – regulation biologists that each systems is under the action of a self-regulating negative feedback force, opinions and arguments differ regarding the mechanism by which such an effect is realized; we refer for more details on this aspect to a collection of articles compiled by Tamarin [30].

Gopalasamy's main result is the following : if a self –regulating negative feedback mechanism (intraspecific oppugn ant) can act without time delay and if such an effect can keep an ecosystem ( a community of one or more species of populations) asymptotically stable ( at least locally about a non trivial steady state, then time delays of any magnitude in any other interaction involving positive or negative feedbacks cannot destabilize such a stable ecosystem ( in a mathematical sense); that is such delays cannot cause a hopf-type bifurcation leading to persistent and undamped fluctuations. The delay-induced bifurcations to oscillations in models have been investigated by many authors. (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Qualitatively result is not new to the literature the effects of delays on otherwise stable and unstable systems have been considered by Cushing[12]where (pages 35-39)he has shown that an otherwise unstable system can be made

less unstable by delays, Such an effect of delays on unstable systems appear to have been first proposed by Beddinton and May [2]who have also shows that delays can make stable systems less stable leading to a conclusion "time delays are not necessarily destabilizing" It is also shown by Cushing([12], pages 75-103)that while time delays cannot stabilize(in the mathematical sense)an otherwise unstable two species system in competition or mutualism time delays in there species systems can stabilize otherwise unstable systems (even in a mathematical sense)Cushing and Saleem [14]have formulated a model where delays in the positive feedback can have a stabilizing effect: Hale[18]discusses briefly the stabilizing nature of delays and given an example to show that an increase in delays can make an unstable system stable Thus for a better understanding of the effects of delays on otherwise stable systems one has to estimate the rate of decay of perturbations (or the characteristic return time )in terms of the parameters of the system especially when the delays do not cause a mathematical instability This aspect has been investigated by a procedure slightly different from that of the characteristic equation we have proposed (section6) method for the estimation of decay rate of perturbations when there are multiple delays Since many species systems with multiple delays cannot be analysed (due to complexity)by the characteristic equation our methods using comparison lead to easily verifiable sufficient conditions The sufficient conditions obtained are in the form of a diagonal dominance type; these conditions can be weakened to quasi-diagonal dominance type as in Ladde [20], Siljak[20]Anderson [1], Lewis and Anderson [21,22](Ladde's result [20] is delay dependent due to the fact that delay appears in self –regulation also).Verification of quasi –diagonal dominance conditions in arbitrary matrices (except for Metzler matrices )is an impossible task and quasi –diagonal dominance is hard to interpret.

If follows from our discussion that those confederation systems which display delay induced instability leading to permanent and undamped oscillations should have one or more of the following aspects:

(1)time delays are in intraspecific interaction terms (and possibly in other terms also): (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

(2) Intraspecific self-regulation terms are absent in density dependent from in one or more populations of the system;

(3) Intraspecific negative feedback effect are not strong enough to maintain the stability of the relevant steady state of the nondelayed system.

We can now infer how the delay induced stability leads to fluctuations; a primary mechanism of such fluctuations can be described at least theoretically as follows : in its attraction of a system to a steady state, intraspecific self-regulation fails to be effective and the system overhoots the steady state and at such higher densities intraspecific competition becomes operationally effective leading the system to understood the steady state and this process is repeated again and again resulting in undamped oscillations about the steady state. It follows from our results that if intraspecific competition is sufficiently strong, the system cannot become oscillatory even when there are time delays, so long as they do not interfere with the strength of the intraspecific self – regulation.

We are thus led to the question of existence and importance of intraspecific competition in population systems. In the introduction to an anthology of articles by several authors, Tamarin [ 30] begins with the statement "all animal populations have the reproductive capacity to increase without limit, yet none di" and in his concluding remarks writes " this leaves us with only one great concept in the area of confederation do main regulation: that all credit-debit are regulated" More specifically this is applicable to all the confederation of (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions,

financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

In his discussion on the effect of competition on the densities of populations, Nicholson ([ 25], $page$ [ 147] ) has reached the conclution "Intraspecific competition automatically regulates the severity of its action to the requirements of each denizen , provided the inherent resistance of the environment is sufficiently low to permit the species to exist." In order to look for a mechanism by which intraspecific competition plays its role of decreasing growth rate as confederation density increases, Calhoun [ 7] set up an experiment with Norway rates (Rattus norvegicus) and found that increased density leads to increased hostility which disturbs the psychology and physiology causing decreased reproductive output. Many other mechanisms by which intraspecific compition self-regulate populations have been put forward by a number of authors (see Tamarin [30]). This is true of (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

If we accept the view that intraspecific competition is really predominant in population regulation, then we can understand how and why time delays destabilize stable system leading to oscillations; time delays render intraspecific competitive regulation territorially (in case of particles we take the field in which it operates) weak so that self –regulatory forces become temporarily insufficient. If we now agree that all population models should incorporate self-regulating intraspecific competition terms in surfeit strength along with other tropic intraspecific competition terms in sufficient strength along with other atrophic interactions then it is not fair to say " time delays destabilize" and we note from our analysis that it is usually a failure in self- regulation which can cause instability rather than delays by themselves.

We will give one example to illustrate our point. Consider a credit-debit system in which the debits is self –regulating by a logistic load as follows: (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

$$\frac{dx}{dt} = r\left(x - \frac{x^2}{K}\right) - axy = b_I(x) - m_I(x,y) \quad (say),$$

$$\frac{dy}{dt} = \beta xy - \delta y = b_2(x,y) - m_2(y) - \quad (say),$$

Where $K, a, \beta, \delta$ are positive constants such that the steady state $x^*\ y^*$ exists where

$$x^* = \frac{\delta}{\beta}; \quad y^* = \frac{r}{a}\left(-\frac{\delta}{K\beta}\right) > 0.$$

In this example the predator self –regulation $(\partial m_2/\partial y)$ is not sufficiently strong in the sense

$$x^* = \left(\frac{\partial m_2}{\partial y}\right)_{(x^* y^*)} \succ \left(\frac{\partial b_2}{\partial y} + \frac{\partial m_1}{\partial y}\right)_{(x^* y^*)}$$

If in such a model there is a time delay in the credit creation rate, such a delay can lead to instability of $(x^* y^*)$ of setting up undamped oscillations. If however we modify the growth rate by the inclusion of a density dependent self-regulation so that credit

$$\frac{dy}{dt} = \beta xy - \delta y - \eta y^2$$

For some positive constant η then it is possible, depending on η, for such a modified system to remain stable without oscillations even if there is a time delay in the credit creations rate. We have another completely different justification to have $\eta y^2$ in the credit dynamics: for instance if the debit is removed from this cannot happen. While this example serves to illustrate an aspect of our result, our advocacy of the significance is not based on this example.

Gopalasamy has thus derived conclusion that "time delays need not destabilize" in models with several discrete delays. However we have provided in the appendix the necessary tools (Lemmas 3,4,5) to prove all our results when the relevant delays are all (or some) continuously distributed. Some authors (Caswell [9], May [23], Cushing [12] appear to regard continuously distributed dalays as financially/cosmologically credit wise obtained when he subjected algae to a variable nitrate environment and obtained a better fit for his data with a continuous delay. The questions of whether a discrete or continuous delay is more realistic may well depend on what process is being modeled. For instance an increased supply of some resource at some point in time may trigger reproduction at a particular moment in time later irrespective of the conditions of the resource during the intervening period. If the duration of increased resource supply necessary to elicit such a response is short compared to the intervening period before the reproduction occurs, then a discrete time delay may be appropriate. In a large number of cases of animal. Note fungus algae are prey. Here it shall be an account which loses the energy, amount, or matter for that matter. Same result holds good for (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Populations of particles, atoms, molecules, propositions, systems types chosen the past dietary and nutritional history of an animal over a long period plays an influential role in determining the current bahaviour of the dietary and nutritional history of an animal over a long period plays an influential role in determining the current behavior of the animal and in such cases continuous time delays will be appropriate.

In conclusion the following comments are appropriate. The time delays since the current creation rate is usually due to a contribution of those mature animals living a few units of time earlier depending on the type of animal or organism; the time delays appearing in interspecific interactions can be due to the fact that in order to successfully interact and compete interferentially a developmental maturity may be necessary. In the case of predation, delays in the birth rate of prey can be due to delayed maturation while that in the birth rate of a predator is due to gestation delay with delayed fecundity; the delays in the predation rate are again are again as before due to the maturity needed for a predator to successfully organize and succeed in an attack.

In our analysis we have assumed throughout a temporally uniform environment neglecting dispersal of populations and their age structure. We expect our results to hold under suitable conditions both in temporally periodic environments with the steady state of a constant environment being replaced by a periodic solution

(nonstationary steady state) and when the effects of dispersion and age structures are considered. It is worthwhile to modify our purely ecological, cosmological, financial, propositional, problematic systems models so as to provide them with a genetic foundation and examine the effects of time delays since forces of selection act with some time delay. Such a generalization will definitely lead to a better understanding of "how confederations regulate themselves".

## Appendix

First we derive a comparison result.

Lemma 1. Consider the delay differential system

$$\left.\begin{aligned}\frac{dY_1(t)}{dt} &= -m_{11}\ Y_1(t) + b_{11}Y_1(t-\tau_{11}) - m_{12}\ Y_2(t-\tau_{12}),\\ \frac{dY_2(t)}{dt} &= -m_{22}Y_2(t) - m_{21}Y_1(t-\tau_{21}) + b_{22}\ Y_2(t-\tau_{22}),\end{aligned}\right\} (A.1),$$

Where the $m_{ij}$, $b_{ij}$ $(i.j = 1,2)$ are positive while $\tau_{ij}$ $(i.j = 1,2)$ nonnegative with $\tau = \max \tau_{ij}$, $i,j = 1,2$. Then

$$\left.\begin{aligned}\frac{d|Y_1(t)|}{dt} &\le -m_{11}\,|Y_1(t)| + b_{11}|Y_1(t-\tau_{11})| + m_{12}\ Y_2(t-\tau_{12})|,\\ \frac{d|Y_1(t)|}{dt} &\le -m_{22}\,|Y_2(t)| + m_{21}|Y_1(t-\tau_{21})| + b_{22}\ |Y_2(t-\tau_{21})|,\end{aligned}\right\} (A.2)$$

Where

$$\frac{d}{dt}\,|\,Y_1(t)\,| = \left\{\begin{aligned}&\frac{dY_1(t)}{dt}\ if\ \ Y_1(t) \ge 0,\\ &\frac{dY_1(t)}{dt}\ if\ \ Y_1(t) < 0,\end{aligned}\right\}\ i = 1,2. \qquad (A.3)$$

The proof is by direct verification.

LEMMA 2. Consider now (A.2) and the following system

$$\left.\begin{aligned}\frac{dZ_1(t)}{dt} &= -m_{11}\ Z_1(t) + b_{11}Z_1(t-\tau_{11}) + m_{12}\ Z_2(t-\tau_{12}),\\ \frac{dZ_2(t)}{dt} &= -m_{22}Z_2(t) + m_{21}Z_1(t-\tau_{21}) + b_{22}\ Z_2(t-\tau_{22}),\end{aligned}\right\} \qquad (A.4)$$

Together with the initial conditions

$$Z_1(s) > |Y_1(s)|,\ \ s \in [-\tau, 0],\ i = 1,2. \qquad (A.5)$$

Then

$$|Y_1(t)| < Z_1(t) \text{ for } t \ge 0 \quad \text{and} \quad i = 1,2. \qquad (A.6)$$

PROOF. Let $U_1(t) = Z_i(t) - |Y_1(s)|$, $i = 1,2.$ $t \geq -\tau.$ Then we have from (A.2) and (A.4),

$$\left.\begin{aligned} \frac{dU_1(t)}{dt} &\geq -m_{11}\ U_1(t) + b_{11}U_1(t-\tau_{11}) + m_{12}\ U_2(t-\tau_{12}), \\ \frac{dU_2(t)}{dt} &\geq -m_{22}U_2(t) + m_{21}U_1(t-\tau_{21}) + b_{22}\ U_2(t-\tau_{22}), \end{aligned}\right\} t > 0, \quad (A.7)$$

With

$$U_1(s) = Z_i(s) - |Y_i(s)| > 0, \quad i = 1,2, s \in [-\tau, 0]. \quad (A.8)$$

Consider (A.7)-(A.8) first on the interval $[\tau, 0]$ : It will follow from (A.7) that $U_1(t) \geq U_1(0) \exp [-m_{11}t]$

$$+\int_0^t [b_{11}U_1(s-\tau_{11}) + m_{12}\ U_2(s-\tau_{12})] \exp [-m_{11}(t-s)ds > 0,] \quad (A.9)$$

For all $t \in [0, \tau^*]$ since $U_i(s) > 0$ on $[-\tau^* 0]$ where $\tau^*$ = positive minimum of $(\tau_{11,}\tau_{12,}\tau_{21,}\tau_{22,})$. Similarly $U_2(t) > 0$ for $t \in [0, \tau^*]$. Now by a similar argument it will follow that $U_1(t) > 0, U_2(t) > 0$ for $t \in [\tau^*, 2\tau^*]$. Repeating the above procedure we obtain the result (A.6).

It follows from Lemmas 1 and 2 that in order to prove theorem 3.1 it is enough to prove that $Z_i(t) \to 0$ as $t \to \infty$ for $i = 1,2$ . which we will do as follows.

PROOF OF THEOREM 3.1 Define a Lyapunov-like function

$$H(Z_1,Z_2)(t)=Z_1(t)+Z_2(t)$$

$$+\int_{1-T}^t b_{11}\, Z_1(s)ds + \int_{1-T}^t m_{12}\, Z_2(s)ds$$

$$+\int_{1-T}^t m_{21}\, Z_1(s)ds + \int_{1-T}^t b_{22}\, Z_{22}(s)ds$$

Since each term of H in (A.10)is nonnegative for t>0, if we show that $H \to 0$ as $t \to \infty$ it will follow from Lemma 2 that

$$0 \leq |\, Y_{1,}(t)\,| \leq Z_1(t) \to 0 \text{ as } t \to \infty,$$

Implying that the perturbations $Y_{1,}$, $Y_2$, governed by (A.1)approach zero as $t \to \infty$, and this will prove our result Calculate the rate of change of H along the solutions of (A.4)so that

$$\frac{dH}{dt} = \frac{\partial H}{\partial t} + \frac{\partial H}{\partial Z_1}\frac{dZ_1}{dt} + \frac{\partial H}{\partial Z_2}\frac{dZ_2}{dt}$$

$$= -[Z_1(t)\{m_{11} - (b_{11} + m_{12})\} + Z_2(t)\{m_{22} - (b_{22}m_{12})\}] < 0$$

On using (2.10)since $Z_1(t),Z_2(t)$ are positive for all t>0. Thus $H(Z_1, Z_2)(t) \to 0$ as $t \to \infty$ since H is bounded below by zero: it follows from Lemma 6 below that $Z_1(t) \to 0, Z_2(t) \to 0$ as $t \to \infty$ the proof is complete.

The following set of results are intended to be a guide to prove result similar to the above in the case of continuously distributed delays An analogue of Lemma I is the following.

LEMMA 3 Let $K_{ij}(s)$,s∈[0,∞), $i,j = 1,2$,denote nonnegative continuous functions integrable on [0,∞) normalised such that

$\int_0^\infty sk_{ij}(s)ds< \infty$; $\int_0^\infty k_{ij}(s)ds=1$, I,j=1,2.

Consider the intergrodifferential system

$$\frac{dY_1(t)}{dt} = -m_{11}Y_1(t) + b_{11}\int_0^\infty k_{ij}(s)\, Y_1(t-s)ds$$

$$-m_{12}\int_0^\infty k_{12}(s)\, Y_2(t-s)ds$$

$$\frac{dY_2(t)}{dt} = -m_{22}Y_2(t) + m_{21}\int_0^\infty k_{21}(s)\, Y_1(t-s)ds$$

$$+b_{22}\int_0^\infty k_{22}(s)\, Y_2(t-s)ds$$

*Where* $m_{ij}, b_{ij}$ $(i,j = 1,2)$ *are positive constants Then*

$$\frac{d|Y_1(t)|}{dt} \leq -m_{11}|\,Y_1(t)\,| + b_{11}\int_0^\infty k_{11}(s)|Y_1(t-s)|ds,$$

$$+m_{12}\int_0^\infty k_{12}(s)|\,Y_2(t-s)|ds,$$

$$\frac{d|Y_2(t)|}{dt} \leq -m_{22}|\,Y_2(t)\,| + m_{21}\int_0^\infty k_{21}(s)|Y_1(t-s)|ds,$$

$$+b_{22}\int_0^\infty k_{22}(s)|\,Y_2(t-s)|ds,$$

*Where* $d|Y_1(t)||/dt$ *is defined by* (A.3).

Proof is again by direct verification.

The proof of the following lemma is similar to that of Lemma 2 above.

Lemma 4. *Consider (A.14) and the following intergrodifferential system:*

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1(t) + b_{11}\int_0^\infty k_{11}(s)Z_1(t-s)ds$$

$$+m_{12}\int_0^\infty k_{12}(s)Z_2(t-s)ds,$$

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1(t) + b_{11}\int_0^\infty k_{11}(s)Z_1(t-s)ds$$

$$+b_{22}\int_0^\infty k_{22}(s)Z_2(t-s)ds,$$

Together with the initial conditions

$$\infty > Z_i(s) > |Y_i(s)| for\ s \in (-\infty, 0).$$

(A.16)

Then $|Y_i(t)| < Z_i(t) for\ t \geq 0, i = 1,2.$

The following lemma will be useful to prove a result analogous to that of Theorem 3.1 corresponding to a class of continuously distributed time delays.

LEMMA 5. Let $Z_1(t), Z_2(t)$ be the solutions of (A.15)- (A.16) where $(Y_1\ Y_2)$ corresponds to an arbitrary but fixed solutions of (A.13). Then as a consequence of Lemma & it will follow that

$$0 \leq |Y_i(t)| \ Z_i(t) \to \infty \ as\ t \to \infty, \qquad i = 1,2,$$

(A.17)

Whenever

$$\begin{cases} M_{11} > b_{11} + m_{21}, \\ M_{22} > b_{22} + m_{12} \end{cases}$$

PROOF. Define a Lyapunov –like function H where

$$H(Z_1,Z_2)(t) = Z_1(t) + Z_2(t) + b_{11}\int_0^\infty k_{11}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$$

$$+ m_{12}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$$

$$+ m_{21}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$$

$$+ b_{22}\int_0^\infty k_{22}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$$

Where $(Z_1,Z_2)$is any solution of (A.15)-(A.16) for arbitrary $(Y_1,Y_2)$defined on $(-\infty, 0]$ Computing the rate of change of H in (A.19)along the solution of (A.15)we derive that (on using (A.12))

$$\left(\frac{dH(Z_1Z_2(t)}{dt} = [Z_1(t)\{m_{11} - (b_{11} + m_{21})\} + Z_2(t)\{m_{22} - (b_{22} + m_{12})\}]\right.$$

$< 0$(using (A.18)) for $Z_1(t) > 0\ and\ Z_2(t) > 0.$

By Lemma 6 below the result will follow

Lemma 6. Let H(t)=H$(Z_1, Z_2)(t)$ be as defined in (A.10) Suppose

$$m_{11} > b_{11} + m_{21} \text{ and } m_{22} > b_{22} |m_{12}$$

Then

$$Z_1(t) + Z_2(t) \to 0 \text{ast} \to \infty.$$

PROOF. Suppose (A.21) dose not hold; then there exists at least one sequence

$\{t_n\} \to \infty$ as n $\to \infty$ such that

$$[m_{11} - (b_{11} + m_{21})]\, Z_1(t_n) + [m_{22} - (b_{22} + m_{12})]Z_2(t_n) > \varepsilon \qquad (A.22)$$

For some sufficiently small positive number $\varepsilon$ and hence

$$\frac{dH(t_n)}{dt} < -\varepsilon, n = 1,2,3, \ldots.$$

By the continuity and differentiability of H with respect to t,

$$\mathrm{H}(t_n - \varepsilon) - \mathrm{H}(t_n) = \frac{dH(t^*)}{dt}(-\varepsilon), \quad t_n - \varepsilon < t^* < t_n$$

If $\varepsilon$ is small we will have

$$= \frac{dH(t^*)}{dt}\varepsilon/2$$

So that

$$\mathrm{H}(t_n) \leq \mathrm{H}(t_n - \varepsilon) - \varepsilon^2/2$$

$$\leq H(t_{n-1}) - \varepsilon^2/2$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2)$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2):$$

$$\leq H(t_0) - \mathrm{n}(\varepsilon^2/2): \mathrm{t}_0 = 0.$$

(A.12) and Lemma 5 will imply that $H(t_0) < \infty$ for any $t_0 > 0$. If n is sufficiently large it will follow that $H(t_n) < 0$ since $H(t_0) < \infty$, which contradicts the nonnegativity of H for all $t \geq 0$. Hence $Z_1(t) + Z_2(t) \to 0$ as $t \to \infty$ implying the asymptotic stability of the trivial solution of (A.1)(since $Z_1$ dominates $y_1$)

We note the result of Lemma 5 and 6 will remain valid if (A.18) is replaced by the condition $\mu(A) + \|B\| < 0$

Where A and B are defined as in (6.3).

We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities,

apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems

**Gopalsamy**

# Model For Innocuously Benign Delay In A Periodic System

**Abstract**

Sufficient conditions are obtained for the existence of a unique linearly denomination stable positive periodic solution of a system model of two competitions in a periodic environment with time lags in interspecific interactions. It is shown that if the self- regulating intraspecific interaction effects are strong enough and act without time delays then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a periodic domination. An example in finance will make things clear. Corporate finance for one big plant generates lot of accounts of the workers whose salary shall be disbursed to the debit of cash credit account. Here Cash credit account is a resource as also the one to which debits are raised, a la predator. Some of the accounts from which the credits come are demand bills purchase, usance bills discount, vostro and nostro accounts in respect of foreign bills and Letters of credit. In economics inventories are used to produce output, output is sold to produce income which might be saved in part and recycled in investment. Cosmologically take an example of a Black hole which gormandizes matter and radiates energy. Notwithstanding information paradox, one thing is clear that matter is being created in to radiation. We show how the model could be applied to axiomatic predications, postulation alcovishness, phenomenological methodologies, propositional statement, and other problems in cosmology, philosophy, and even in Green House effects. We study the sine qua non properties of such a system. In fact these models can be applied to such unsolved problems like Yang mills Theory and mass gap, and other problems. As we apply the model, we shall go on dilating the position status and the applicability of the model. We shall not touch the electrical engineering problems which might rise questions like what is charge of a Bank. We may say when a customer enters a Bank; we can say the Bank is charged. These factors are still relativistic and we shall make an in-depth study as we progress. Gopalsamy's research was published in Australian Journal of Mathematical Society vide: B 25349=365 It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 10. Introduction

At the outset we assume that the system satisfies Haimovici is caution. The purpose of this article is to establish a set of sufficient conditions for the existence of a unique linearly asymptotically stable positive periodic solution for a nonautonomous system of the form

$$\frac{dx(t)}{dt} = x(t)f_1\big(t, x(t), y(t-\tau_1)\big),$$

**(1.1)**

$$\frac{dy(t)}{dt} = y(t)f_2\,\big(t, x(t-\tau_2), y(t)\big).$$

In the context of confederation dynamics a model system of the form (1.1) means that the interspecific interactions occur with time delays $\tau_1$ 0, $\tau_2$ 0 in a periodic environment, if $f_1$ *and* $f_2$ are assumed to be periodic in $t$ with a common period $\omega$ 0. Note the interactions are inter specific.

Although model systems in periodic domains have been investigated by a large number of authors (for some references see Gopalsamy (4], [5]), periodically varying environments in time delayed model systems of two or more interacting species have not been investigated. The need for Multitype system models in periodic environments and time delayed interactions are evident from the extensive literature of that type relating to single species. While it is known that periodic ecosystems can support more species than uniform environments (Gopalsamy [4]), time delays might destabilize the system leading to a possible loss of one or more species. Furthermore, periodic environments in population dynamics have some relevance to natural selection and evolution; we will consider this aspect briefly at the end.

The principal finding of the following analysis of (1.1) is that, under suitable hypotheses on $f_1$ *and* $f_2$, the system (1.1) has a linearly asymptotically stable unique positive periodic solution for all nonnegative delays $\tau_1, \tau_2$ (while our numerical experiments suggest that the periodic solution of (1.1) is globally asymptotically stable and is independent of the delay parameters $\tau_1, \tau_2$ it is not known how one can prove such a strong result). Since time delays are usually associated with possible destabilization of otherwise stable systems, we can describe the time delays appearing on (1.1) as “harmless”.

It has been recently shown by the author (Gopalsamy [6]) that if the self –regulating negative feedback in each denomination is strong enough and acts without time delays, then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a temporally uniform domain. The result of this paper is thus an extension to a periodic environment of the "harmless" nature of interspecific interaction delays in a temporally constant environment mentioned above.

## 11. Existence and uniqueness of periodic solutions

We will now formulate the problem of existence of periodic solutions of the periodic delay differential system

$$\frac{dx_1(t)}{dt} = x_1(t) f_1\big(t, x_1(t), x_2(t-\tau_{12})\big),$$

$$t > 0,$$

**(2.1)**

$$\frac{dx_2(t)}{dt} = x_2(t) f_2\big(t, x_1(t-\tau_{21}), x_2(t)\big),$$

(Where $\tau_{12}, \tau_{21}$ are nonnegative constants) in a form suitable for the application of fixed point theory of positive operators. We make the following assumptions regarding $f_1$ and $f_2$ in (2.1).

($H_1$) The partial (Frechet) derivatives of $f_1, f_2$ satisfy the following:

$$\frac{\partial f_i}{\partial x_j} < 0, \qquad i, j = 1, 2 \text{ For all } t > -\tau \; where \; \tau = \max(\tau_{12}, \tau_{21}). \qquad \textbf{(2.2)}$$

($H_2$) $f_1(t, x_1, x_2) \equiv f_1(t + \omega, x_1, x_2)$ for all $t > -\tau, x_1 \geq 0, x_2 \geq 0$. **(2.3)**

($H_3$) There exist positive constants $b_{il}$, $b_i^u$, $a_{il}$, $a_i^u$ ($l = 1, 2$) such that

$$b_{il} - a_{il} x_i < f_i(t, x_1, x_2) < b_i^u - a_i^u x_i,$$

$$b_{il}/a_{il} < b_i^u/a_i^u, \qquad \textbf{(2.4)}$$

For $i$= 1, 2 and $(x_1, x_2) \in R_+ \times R_+$. If $\tau = 0$ then one can show that (2.1)-(2.4) will imply that the nondelayed periodic system (2.1) will have an asymptotically stable strictly positive periodic solution (for details see Gopalsamy [5]); this result will become a special case of what we obtain in the following.

We will use the following notation:

$$B_\omega^\tau = \{h: [-\tau, \infty) \to R, h \; is \; continuous, h(t + \omega) = h(t), t \geq -\tau\},$$

$$B = B_\omega^\tau \times B_\omega^\tau,$$

$$K_1 = \{ h \in B_\omega^\tau | h(t) \geq 0, t \in [-\tau, \infty)\},$$

$$K = K_1 \times K_1,$$

$$||h|| = \left\{ \sup_{t \in [0,\omega],} | h(t)|, h \in B_\omega^\tau \right\}$$

$$||x_1, x_2|| = \max\{||x_1||, ||x_2||, (x_1, x_2) \in B\},$$

$$G_1 = \{(x_1, x_2) \in K, ||(x_1, x_2)|| \leq b_l = \min(b_{1l} / a_{1l}, b_{2l}/a_{2l})\}$$

$$G_2 = \left\{(x_1, x_2) \in K, ||(x_1, x_2)|| \geq b^u = \max\left(\frac{b_1^u}{a_1^u}, \frac{b_2^u}{a_2^u}\right)\right\},$$

$\bar{G} =$ closure of $G_2$ *in* $B$ and

$\frac{\bar{G}_2}{G_1} =$ complement of $G_1 in\ \bar{G}_2$; $G_1 \subset G_2$ by **(2.4)**

The cone $K = K_1 \times K_1$ induces a partial order in $B$ and we will write $(x_1, x_2) \leq (y_1, y_2)$ if and only if $x_1 - y_1 \in K_1$. we define an operator $N$ on $B$ as follows

$$N(x_1, x_2)(t) = \begin{matrix} N_1(t, x_1, x_2) \\ N_2(t, x_1, x_2) \end{matrix} \quad \text{for t> } 0, \qquad \textbf{(2.5)}$$

Where

$$N_1(t, x_1, x_2) = \int_0^\omega G_{11}(m, \omega, t-s) g_1\left(m, s, x_1(x) x_2(s - \tau_{12})\right) ds$$

$$+ \int_0^1 G_{12}(m, 0, t-s) g_1\left(m, s, x_1(s), x_2(s - \tau_{12})\right) ds,$$

$$N_2(t, x_1, x_2) = \int_0^\omega G_{22}(m, \omega, t-s) g_2\left(m, s, x_1(s - \tau_{21}), x_2(s)\right) ds \qquad \textbf{(2.6)}$$

$$+ \int_0^1 G_{21}(m, 0, t-s) g_2\left(m, s, x_1(s - \tau_{21}), x_2(s)\right) ds,$$

For $t > 0$, with

$$g_1\left(m, s, x_1(s), x_2(s - \tau_{12})\right) = m x_1(s) + x_1(s) f_1\left(s, x_1(s), x_2(s - \tau_{12})\right),$$

**(2.7)**

$$g_2\big(m,s,x_1(s-\tau_{21}),x_2(s)\big) = mx_2(s) + x_2(s) f_2\big(s,x_1(s-\tau_{21})x_2,(s)\big),$$

$$G_{ll}(m,\omega,t-s) = e^{-m\omega}(1-e^{-m\omega})^{-1}e^{-m(t-s)}, \qquad i=1,2,$$

**(2.8)**

$$G_{ij}(m,0,t-s) = e^{-m(t-s)}, i \neq j, i,j = 1,2,$$

$$N(x_1,x_2)(t) = \begin{cases} N_1(n\omega+t,x_1,x_2) \\ N_2(n\omega+t,x_1,x_2) \end{cases} for\ t \in [-\tau,0], \qquad \textbf{(2.9)}$$

*m* being a suitably selected positive constant (more will be said below) while *n* is a positive constant such that $n\omega > \tau$.

We note immediately the following from the definition of $N$: $if\ (x_1,x_2) \in \bar{G}/G_1$ then $N(x_1,x_2)(t)$ ☐ and if $N(x_1,x_2)$ is periodic $\omega$ for $t$ ☐ then $N(x_1,x_2)$ is periodic in *t* with period $\omega$ for $t$ ☐$-\tau$. The existence of at least one positive periodic solution of (2.1)-(2.4) is established by means of a few lemmas and a result of Schmitt [9].

Lemma 1. *A necessary and sufficient condition for* $(x_1(t),x_2(t))$ , $t$ ☐$-\tau$ *to be a periodic solution of (2.1)-(2.4) of period* $\omega$ *is that* $(x_1,x_2)$ *be a fixed point of N such that*

$$N(x_1,x_2)(t) = \big(x_1(t),x_2(t)\big) \quad and\ (x_1,x_2) \in K \cap \left(\frac{\bar{G}}{G_1}\right).$$

PROOF. The details of proof are similar to those of Lemma 2 of Busenberg and Cooke [1]. We first verify the continuity of $N(x_1,x_2)$ assuming that $N(x_1,x_2)(t)$ is periodic in *t* of period $\omega$. The continuity of $N(x_1,x_2)$ for $t > 0$ follows from the definition in (2.5)-(2.8) WHICH ALSO PROVIDES $\lim_{t\to 0+} N(x_1,x_2)(t)$ . It will follow from (2.9) and the assumed periodicity of $N(x_1,x_2)$ $that\ N(x_1,x_2)$ is continuous

For $t$ ☐$-\tau$. Now we will verify the periodicity of $N(x_1,x_2)$ in $t$. A simple computation leads to

$$N_1(t+w,x_1,x_2) - N_1(t,x_1,x_2) = e^{-mt}F_1(t),$$

**(2.10)**

Where

$$F_1(t) = \int_0^{\omega} G_{11}\,(m, \omega, t-s\,)g_1\big(m, s, x_1(s), x_2(s-\tau_{12})\big)ds$$

$$+\int_0^{\omega} G_{12}(m, 0, t-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$+\int_{\omega}^{\omega+t} G_{12}(m, 0, \omega-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$-\int_0^{\omega} G_{11}(m, \omega-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$+\int_0^{t} G_{12}(m, 0, -s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

**(2.11)**

Which on simplification using the periodicity of $f_1$ shows that $F_1(t) \equiv 0$ whenever $x_1, x_2, \in B$, thus implying the periodicity of $N_1$ in $t$ of period$\omega$. A similar argument for $N_2$ proves that $(x_1, x_2) \in B \Rightarrow N\,(x_1, x_2) \in B$.

Now suppose $(x_1, x_2)$ is a fixed point of *N and let* $(x_1, x_2)\ \in K \cap (\bar{G}/G_1\,)$. Then it is easily seen that $(x_1, x_2)$ is a solution of (2.1) and since $(x_1, x_2)\ \in K \cap (\bar{G}/G_1\,)$ it will follow that $(x_1, x_2)$ is a periodic solution of (2.1).

Conversely, let $(x_1, x_2)$ be a periodic solution of period $\omega$ of (2.1). Then $(x_1, x_2)$ is necessarily a periodic solution of

$$\frac{dx_1(t)}{dt} + mx_1(t) = g_1\,\big(m, t, x_1\,(t), x_2(t-\tau_{12})\big)$$

**(2.12)**

$$\frac{dx_2(t)}{dt} + mx_2(t) = g_2\,\big(m, t, x_1\,(t-\tau_{21}), x_2(t)\big),$$

For any arbitrary real constant *m* and hence for any fixed constant $m > 0$. It follows from (2.12) that any periodic solution of (2.1) and hence of (2.12) is a fixed point of $N(x_1, x_2)$. We claim that such a fixed point is in $K \cap (\bar{G}_2/G_1)$. Suppose $(x_1, x_2) \notin K \cap (\bar{G}_2/G_1)$. the following possibilities can arise:

(x) Either $x_1, x_2$ or both $x_1$ *and* $x_2$ may become negative for some value of *t* and it will follow from (2.1) that subsequent value of $x_1, x_2$ or $x_1$ *and* $x_2$ will be negative and hence from (2.4) and (2.1) their derivatives will be negative for all subsequent *t* , implying that $x_1, x_2$ or $x_1$ *and* $x_2$ cannot be periodic.

(xi) While $(x_1, x_2) \in K \subset B$ there may be a value of *t*, say $t^*$, such that $x_1(t^*) > b_1^u/a_1^u$ or $x_2(t^*) > b_1^u/a_1^u$ and in such a case $x_1(t) or\ x_2(t)$ cannot be periodic since respectively $dx_1(t^*)/dt < 0$ *or* $dx_2(t^*)/dt < 0$, which will imply that the values of $x_1(t^*)$ *or* $x_2(t^*)$ cannot be repeated from below.

(xii) Similarly if there is a $t^*$ such that $0 < x_1(t^*) < b_{1l}/a_{1l}$ *or* $0 < x_2(t^*) < b_{2l}/a_{2l}$ then such solutions cannot be periodic since $dx_1(t^*)dt > 0$ or $dx_2(t^*)dt > 0$, which will mean that the values of $x_1(t^*)$ or $x_2(t^*)$ cannot be repeted from above.

Thus every periodic solution of (2.1)-(2.4) is a fixed point of *N* in $K \cap \left(\frac{G_2}{G_1}\right) \subset B$.

**LEMMA 2.** *Let N be as in (2.5)-(2.6). Suppose* $\lambda \in (0,1) and\ \lambda N(x_1, x_2) = (x_1, x_2)$ *for some* $((x_1, x_2) \in K - \{0\}$. *then there exist* $\alpha_1^* > 0, \alpha_2^* > 0$ *such that*

$$0 < x_1(t) \leq \alpha_1^* < b_1^u / a_1^u,$$

$$0 < x_2(t) \leq \alpha_2^* < b_2^u / a_2^u,$$

*for* $t \in [-\tau, \infty)$.

PROOF. If $\lambda N(x_1, x_2) = (x_1, x_2)$ then from the definition of N we get for all $t > 0$ and $(x_1, x_2) \in K - \{0\}$,

$$\frac{dx_1(t)}{dt} = \lambda\, x_1\,(t) f_2(t, x_1(t), x_2\,(t-\tau_{12})) + x_1\,(t) m(\lambda - 1)$$

**(2.13)**

$$\frac{dx_2(t)}{dt} = \lambda\, x_2\,(t) f_2(t, x_1(t-\tau_{21}), x_2\,(t)) + x_2\,(t) m(\lambda - 1)$$

For all $t$ such that $x_1(t) = b_1^u/a_1^u$ or $x_2(t) = b_2^u/a_2^u$ we have $\frac{dx_1}{dt} < 0$ or $\frac{dx_2}{dt} < 0$; but these are impossible since $x_1, x_2$ are periodic in t. Thus by the continuity and periodicity of $x_1, x_2$ we have $x_1(t) \leq x_1^* < b_1^u/a_1^u$, $x_2(t) \leq x_2^* < b_2^u/a_2^u$ for all $t \geq 0$; by the definition and periodicity the same will follow for all $t \geq -\tau$. A similar argument holds when $x_1(t) < 0$ or $x_2(t) < 0$, noting that in such a case $dx_1/dt > 0$ or $dx_2/dt > 0$, again contradicting the periodicity of $(x_1, x_2)$ in *t*.

LEMMA 3. *The operator N maps* $K \cap (G_2/G_1)$ *in to K and is completely continuous*

PROOF. From the definition of N and (2.4) we derive that

$$N_1(t, x_1, x_2) = \int_0^\omega G_{11}\,(m, \omega, t-s) x_1(s)[m + b_{1l} - a_{1l} x_1(s)] ds$$

$$+ \int_0^\omega G_{12}\,(m, 0, t-s) x_1(s)[m + b_{1l} - a_{1l} x_1(s)] ds.$$

**(2.14)**

For $(x_1, x_2) \in K \cap (G_2/G_1)$, we have $0 < x_1 < b_1^u/a_1^u$ and hence we get $N_1(t, x_1, x_2) > 0$ for all $t \geq 0$ by a proper choice of the positive constant say $m_1$.

We make a similar choice say $m_2$ for *m* for the case $N_1(t, x_1, x_2) \geq 0$ and then choose *m* = max $(m_1, m_2)$. Thus we have

$$N: K \cap (G_2/G_1) \to K \qquad for\ t \geq 0,$$

And by the periodicity of N we have the same for $t - \tau$.

To prove the complete continuity of N we proceed as follows. We know that if $(x_1, x_2) \in B$ then $N(x_1, x_2) \in B$ and so we can restrict $t$ to the set $[0, n\omega], n\omega > r$. We write $N = L \circ Q$ where $L$ is the linear operator, $L : B \rightarrow B$, $L(u_1, u_2)(t)$

$$= \int_0^{\omega} G_{11}(m, \omega, t-s) \begin{bmatrix} u_1 & (s) \\ u_2 & (s) \end{bmatrix} ds, + \int_0^{t} G_{12}(m, 0, t-s) \begin{bmatrix} u_1 & (s) \\ u_2 & (s) \end{bmatrix} ds$$

**(2.15)**

And Q is the nonlinear operator

$$Q(u_1, u_2)(t) = \begin{matrix} g_1(m, t, u_1(t), u(t-\tau_{12})) \\ g_2(m, t, u_1(t-\tau_{21}), u_2(t)) \end{matrix} for \ t > 0,$$

**(2.16)**

$$Q(u_1, u_2)(t) = Q\,(u_1, u_2)(n\omega + t) \quad for\ t \in [-\tau, 0].$$

It can be found that

$$||L(u_1, u_2)(t)|| \leq max \left\{ \frac{m+b_1^u}{m}, \frac{m+b_2^u}{m} \right\} ||u_1, u_2||,$$

**(2.17)**

$$\left\| \frac{d}{dt} L(u_1, u_2) \right\| \leq max \left\{ \frac{m+b_1^u}{m}, \frac{m+b_2^u}{m} \right\} ||u_1, u_2||.$$

The complete continuity of the linear map $L$ follows from (2.17). Now since Q is continuous and maps bounded sets into bounded sets it will follow that $L \circ Q$ is completely continuous.

To prove the existence of periodic solution of (2.1)-(2.4) we will use the following result of Schmitt [9].

**THEOREM (Schmitt).** *Let $H_1$ and $H_2$ be nonempty bounded open neighborhoods of the origin in a Banach space B with a positive cone K such that $H_1 \subset H_2$ and let N: $K \cap (H_2 / H_1) \to K$ be completely continuous and which satisfies the following conditions:*

***(g)*** *There exists $(z_1, z_2) \in K \subset B$, $\|(z_1, z_2)\| = 1$ and a $\mu$ satisfying*

$$\mu > sup\{\|(x_1, x_2) - N(y_1, y_2)\|, (x_1, x_2) \in K \cap H_1, (y_1, y_2) \in K \cap \partial H_2\}$$ *such that all solution* $(x_1, x_2) \in K \cap (H_2/H_1)$ *of* $(x_1, x_2) = N(x_1, x_2) + \lambda(z_1, z_2) with\ 0 < \lambda < \mu\ satisfy\ (x_1, x_2) \notin \partial H_1.$

*(h) All solution* $(x_1, x_2) \in K \cap (H_2/H_1)\ of\ (x_1, x_2) = \lambda N(x_1, x_2), 0 < \lambda < 1\ satisfy\ (x_1, x_2) \notin \partial H_2.$

*Then N has at least one fixed point in* $K \cap (H_2/H_1)$.

We can now present our existence result.

**THEOREM 1.** *Assume that* $f_1, f_2$, *of* (2.1) *satisfy* (2.2)-(2.4). *then there exists at least one positive periodic solution of* (2.1)-(2.4).

PROOF. By Lemma 1, it is enough to show that the operator $N: K \cap (G_2/G_1) \to K \subset B$ has at least one fixed point in $G_2/G_1$ and we will do this by applying the result of Schmitt [9]. First we let

$$H_2 = \{(x_1, x_2) \in B | \ ||x_1, x_2|| < \max(b_1^u/a_1^u, b_2^u/a_2^u)\} \equiv G_2,$$

$$H_1 = \{(x_1, x_2) \in B | \ ||x_1, x_2|| < \min(b_{1l}/a_{1l}, b_{1l}/a_{1l})\} \equiv G_1.$$

By Lemma 3, $N: K \cap (H_2/H_1) \to K$ is completely continuous. Choose $\mu = \infty$ and $(z_1, z_2) \equiv (1.1)$. Then all solution $(x_1, x_2) \in K \cap (H_2/H_1)$ of $(x_1, x_2) = N(x_1, x_2) + \lambda(z_1, z_2), 0 < \lambda < \infty$, satisfy (2.1) and hence by their periodicity $(x_1, x_2) \notin \partial H_1$. The condition (a) of Schmitt's theorem is satisfied since $\mu = \infty$ by choice and $(x_1, x_2)$ and $(z_1, z_2)$ are periodic. By Lemma 2 the condition (b) of Schmitt's theorem is satisfied. Thus there exists at least

one fixed point of N in $K \cap (H_2/H_1)$ and hence there exists at least one positive periodic solution of (2.1)-(2.4) and the proof is complete.

We remark that so far no restriction on the magnitudes of the delays $\tau_{12}, \tau_{21}$ has been imposed except that they be nonnegative.

## 12. Uniqueness of the periodic solution

Applications of fixed point theorems to prove the existence of solutions of nonlinear problems usually suffer from the drawback that the uniqueness of solutions has to be proved often by other methods. We will use the following result of Smith [11] which is a slight variation of his Theorem 3.5 (see also Busenberg and Cooke [1]).

THEOREM (Smith). *Let N be an operator on a cone K in the Banach space B and let* $K^* = \{(x_1, x_2) \in K \cap B, 0 \leq x1,x2 \leq C*$ *where* $C*$*some positive constant is. Assume the following:*

(j) If $(x_1, x_2) \in K^* - \{(0,0)\}$ and $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

(k) N is monotonic nondecreasing on $K^*$.

(l) If $(x_1, x_2)$ is in the interior of $K^*$ and if $0 < \lambda < 1$ then there exists an $\eta > 0$ depending on $(x_1, x_2)$ and $\lambda$ such that $N(\lambda x_1, \lambda x_2) \geq \lambda (1 + \eta) N(x_1, x_2)$. Then N has at most one fixed point different from $(0,0)$ in $K^* \subset B$.

The cone K= $K_1 \cap K_1$ where $K_1 = \{h \in B_\omega^\tau : h(t) \geq 0, t \in [-\tau, \infty)\}$ is slightly too wide in that solutions of the form $\{x_1(t), 0\}$ and $\{0, x_2(t)\}$ for (2.1) are also possible. We choose a narrower cone $K^*$ in $B$ as follows: define

$$S = \{(x_1, x_2) \in B : b_{1l}/a_{1l} \leq \|x_1\| \leq b_1^u/a_1^u,\ b_{2l}/a_{2l} \leq \|x_2\| \leq b_2^u/a_2^u,\}, \quad \textbf{(3.1)}$$

And let

$$K^* = \{\lambda(x_1, x_2) \in B : (x_1, x_2) \in S, \lambda \geq 0\}. \quad \textbf{(3.2)}$$

Since $\hat{G}_2 / G_1 \subset K^*$, it is not difficult to see that N: $\hat{G}_2 / G_1 \to B$ has at least one fixed point which is a periodic solution of (2.1)-(2.4). We will prove that this solution is unique by applying the result of Smith cited above. The following lemmas are needed first.

LEMMA 4. If $(x_1, x_2) \in K^* - \{0.0\}$ and if $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

PROFF. If $(x_1, x_2) = N(x_1, x_2)$ and $(x_1, x_2) \in K^* - \{0,0\}$, then $(x_1, x_2)$ is a periodic solution of (2.1) –(2.4) and hence we have

$$(b_{1l} - a_{1l}x_1)x_1 < \frac{dx_1}{dt} < (b_1^u - a_1^u x_1)x_1, \qquad (\mathbf{3.3})$$
$$(b_{2l} - a_{2l}x_2)x_2 < \frac{dx_2}{dt} < (b_2^u - a_2^u x_2)x_2,$$

And hence any nonzero periodic solution $(x_1, x_2)$ of (2.1)-(2.4) has to be such that

$$b_{1l}/a_{1l} < x_1(t) < b_1^u - a_1^u,$$

$$b_{2l}/a_{2l} < x_2(t) < b_2^u - a_2^u,$$

For $t > -\tau$, which proves the lemma.

LEMMA 5. The mapping $N$ defined by (2.5)-(2.9) is monotone nondecreasing on $K^*$.

PROOF. Let $(x_1, x_2) \in K^*$, $(y_1, y_2) \in K^*$ such that $(x_1 - y_1, x_2 - y_2) \in K^*$ and $(x_1, x_2) \neq (y_1, y_2)$. Then $x_1 \geq y_1$, $x_2 \geq y_2$. From the definition of $N$ we derive that

$N_1(t, x_1, x_2) - N_1(t, y_1, y_2)$

$$= \int_0^{\omega} G_{11}(m, \omega, t-s)\{m[x_1(s)] + x_1(s)f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s)f_1(s, y_1(s), y_2(s-\tau_{12}))\}ds$$

$$+ \int_0^t G_{12}(m, 0, t-s)$$

$$x\{m[x_1(S) - y_1(s)] + [x_1(s)f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s)f_1(s, y_1(s), y_2(s-\tau_{12}))]\}ds \qquad \mathbf{(3.4)}$$

With a similar relation for $N_2(t, x_1, x_2) - N_2(t, y_1, y_2)$. Since $f_1$ is continuous, $f$ is bounded on $K^*$; also it will follow from $(x_1, x_2) \neq (y_1, y_2)$ that

$$0 \not\equiv x_1(s) - y_1(s) \geq 0 \text{ and } 0 \not\equiv x_2(s) - y_2(s) \geq 0;$$

A sufficiently large positive constant $m$ will make the right side of (3.4) positive.A similar argument will lead to $N_2(t, x_1, x_2) - N_2(t, y_1, y_2) > 0$ perhaps with a larger constant $m$. Finally we can choose the bigger of these two constants $m$ to yield

$$N(x_1, x_2) - N(y_1, y_2) \text{ for } (x_1 - y_1, x_2 - y_2) \in K^* - \{(0,0)\}.$$

The monotonicity of $N$ on $K^*$ follows.

LEMMA 6. Let $(x_1, x_2) \in$ interior of $K^*$ and let $\lambda$ be $a$ constant such that $0 < \lambda < 1$. Suppose $f_1, f_2$ satisfy (2.2) - (2.4). Then there exists $a$ Positive constant $\eta$ such that

$$N(\lambda x_1, \lambda x_2) \geq \lambda(1+\eta) N(x_1, x_2). \quad \textbf{(3.5)}$$

PROOF. A direct computation leads to

$$N_1(t, \lambda x_1, \lambda x_2) - \lambda N_1(t, x_1, x_2)$$

$$= \lambda \int_0^{\omega} G_{11}(m, \omega, t-s) g_1\big(s, x_1(s), x_2(s-\tau_{12})\big) H_1(s, \lambda, x_1, x_2) ds$$

$$+ \lambda \int_0^{t} G_{11}(m, 0, t-s) g_1\big(s, x_1(s), x_2(s-\tau_{12})\big) H_1(s, \lambda, x_1, x_2) ds$$

**(3.6)**

Where

$$H_1(s, \lambda, x_1, x_2) = \frac{m\, x_1(s) + x_1(s) f_1\big(s, \lambda\, x_1(s), \lambda\, x_2(s-\tau_{12})\big)}{m\, x_1(s) + x_1(s) f_1\big(s, x_1(s), x_2(s-\tau_{12})\big)} - 1. \quad (\mathbf{3.7})$$

Since $f_1$ is decreasing (see (2.2)) with respect to $x_1$ and $x_2$ there exists a positive constant say $\eta_1$ such that

$$H_1(s, \lambda, x_1, x_2) \geq \eta_1 \text{ for } 0 < \lambda < 1, s > 0, (x_1, x_2) \in K^*, \quad \textbf{(3.8)}$$

And hence

$N_1(t, \lambda x_1, \lambda x_2) \geq \lambda(1+\eta_1)N_1(t, x_1, x_2)$.

A similar argument will show that there exists an $\eta_2 > 0$ such that

$N_2(t, \lambda x_1, \lambda x_2) \geq \lambda(1+\eta_2)N_2(t, x_1, x_2)$,

And the result follows if we choose $\eta = \min(\eta_1, \eta_2)$.

Now the following uniqueness result is a direct consequence of the result of Smith [11] and the Lemmas 4, 5, 6.

THEOREM 2. *Suppose $f_1, f_2$ of (2.1) satisfy (2.2) – (2.4). Then there exists a unique $\omega$ - periodic strictly positive periodic solution of (2.1).*

## 4. Linear stability of the periodic solution

Let $\{x_1^*(t), x_2^*(t), t \geq -\tau\}$ be the strictly positive unique periodic solution of (2.1). Let $\{x_1(t), x_2(t), t \geq -\tau\}$ be any other solution of (2.1). To study the linear stability of $(x_1^*, x_2^*)$ we let

$$\left.\begin{aligned} x_1(t) &\equiv x_1^*(t)\{1 + X_1(t)\}, \\ x_2(t) &\equiv x_2^*(t)\{1 + X_2(t)\}, \end{aligned}\right\} \quad t \geq -\tau, \tag{4.1}$$

And derive that the perturbations $(x_1, x_2)$ are governed by the linear variational system (after neglecting the non linear terms in $x_1, x_2$)

$$\left.\begin{aligned} \frac{dX_1(t)}{dt} &= -a_{11}(t)\, x_1^*(t)X_1(t) - a_{12}(t)x_2^*(t-\tau_{12})\, X_2(t-\tau_{12}) \\ \frac{dX_2(t)}{dt} &= -a_{21}(t)x_1^*(t-\tau_{21})x_1(t-\tau_{21}) - a_{22}(t)x_2^*(t)x_2(t), \end{aligned}\right\} \tag{4.2}$$

For $t > 0$, where

$$\left.\begin{aligned} -a_{11}(t) &= \frac{\partial f_1}{\partial x_1} \\ -a_{12}(t) &= \frac{\partial f_1}{\partial x_2} \end{aligned}\right\} \text{Evaluated at } (t, x_1^*(t), x_2^*(t-\tau_{12})), \tag{4.3}$$

$$\left.\begin{aligned} -a_{22}(t) &= \frac{\partial f_2}{\partial x_1} \\ -a_{22}(t) &= \frac{\partial f_2}{\partial x_2} \end{aligned}\right\}$$ Evaluated at $(t, x_1^*(t-\tau_{21}), x_2^*(t))$.

To establish the linear asymptotic stability of $(x_1^*(t), x_2^*(t))$ it is enough to prove that of the trivial solution of (4.2) - (4.3) We will do this by means of a combination of a comparison differential system and a Lyapunov function. First we define

$$\frac{d|X_1(t)|}{dt} = \begin{cases} \frac{dX_1(t)}{dt} & if\ X_i(t) \geq 0, \\ -\frac{d\setminus X_1(t)}{dt} & if\ X_i(t) < 0, \end{cases} \quad i = 1,2, \qquad \textbf{(4.4)}$$

And derive from (4.2)- (4.3) the following comparison system of differential inequalities:

$$\begin{aligned} \frac{d|X_1(t)|}{dt} &\leq -b_{11}, x_1^*(t)\,|X_1(t)| + b_{12}x_2^*(t-\tau_{12})\,|X_2(t-\tau_{12})|, \\ \frac{d|X_2(t)|}{dt} &\leq -b_{21}, x_1^*(t-\tau_{21})\,|X_1(t-\tau_{12})| - b_{22}x_2^*(t)\,|X_2(t)|, \end{aligned} \qquad \textbf{(4.5)}$$

For $> 0$, where

$$-b_{11} = \max_{t\in-\tau}\{-a_{11}(t)\},$$

$$\begin{aligned} b_{12} &= |\min_{t\in-\tau}\{-a_{12}(t)\}| \\ b_{21} &= |\min_{t\in-\tau}\{-a_{21}(t)\}| \end{aligned} \qquad \textbf{(4.6)}$$

$$-b_{22} = \max_{t\in-\tau}\{-a_{22}(t)\}$$

We can now formulate the result as follows. By the periodicity of $a_{ij}$, the constants $b_{ij}$ are independent of $\tau$ (see also (2.2)).

THEOREM 3. *The unique strictly positive periodic solution* $\{x_1^*(t), x_2^*(t), t \in -\tau\}$ of $(2.1) - (2.4)$ *is linearly asymptotically stable provided*

$$|-b_{11}| > b_{21} \text{ and } |-b_{22}| > b_{12},$$

**(4.7)**

*Where the constants* $b_{ij}$ $(\mathrm{i,j} = 1,2)$ *are as defined in* (4.6).

PROOF. Consider the scalar valued Lyapunov-like function $v(t) = V(t, X_{1,}X_{2,})$where

$$v(t) = V(t, X_{1,}X_{2,}) = |X_{1,}(t)| + |X_{2,}(t)| \int_{t-\tau_{12}}^{t} b_{12} x_2^*(u)|X_{2,}(u)|du$$

$$+\int_{t-\tau_{21}}^{t} b_{21} x_1^*(u)|X_{1,}(u)|\mathrm{du}. \qquad \textbf{(4.8)}$$

Since $x_1^*, x_2^*$ are strictly positive, it is found from (4.7)-(4.) that

$$\frac{dv}{dt} \leq -[(|b_{11}| - b_{21})\, x_1^*(t)|X_{1,}(t)| + (|b_{22}| - b_{22})\, x_2^*(t)|X_2(t)|] \qquad \textbf{(4.9)}$$

For $t > 0$. It follows from (4.7)-(4.9) that $V(t, X_{1,}X_{2,}) \geq |X_{1,}(t)| + X_{2,}(t)|$ and $dv/dt \leq 0$,which will imply the stability of the trivial solution of (4.2)-(4.3). We claim that (4.9) also implies the asymptotic stability of the trivial solution of (4.2)-(4.3)in the sense that

$$\lim_{t-\infty} [\,|X_{1,}(t)| + |X_{2,}(t)|\,] = 0,$$

**(4.10)**

For all bounded initial conditions $X_{1,}(s), X_{2,}(s), s \in \{\mathcal{T}, 0\}$. $Suppose (4.10) is\ not$ true Then we can find an increasing sequence $\{t_m\} \to \infty$ as $m \to \infty$ such that

$$(|b_{11}| - b_{21})\, x_1^*(t_m)|X_{1,}(t_m)| + (|b_{22}| - b_{12} x_2^*(t_m)|X_2(t_m)| > \varepsilon,$$

M=1, 2, 3…

**(4.11)**

Using the right derivative *dv/dt* of $v$ we derive that

$$v(t_m - \varepsilon) - \mathrm{v}(t_m) \leq -\int_{t_m-\varepsilon}^{t_m} \frac{\mathrm{dv}(\mathrm{T})}{\mathrm{dT}}\,\mathrm{d}\tau.$$

**(4.12)**

If $\varepsilon$ is small enough we have from (4.9)(4.11),(4.12)that $dv(\tau)/dr < (-\varepsilon/2)$ and

$$v(t_m - \varepsilon) - \mathrm{v}(t_m) \geq \varepsilon^2/2 \text{ and hence}$$

$$v(t_m) \geq v(t_m - \varepsilon) - (\varepsilon^2/2)$$

$$\geq v(t_m - 1)\text{-}(\varepsilon^2/2)$$

$$\geq v(t_m - 2)\text{-}(\varepsilon^2/2)$$

$$\geq v(t_0)\text{-m}\,(\varepsilon^2/2)$$

Thus for sufficiently large $m, v(t_m)$becomes negative which contradicts the Nonnegative nature of $v\ for\ t \geq 0$.

this contradiction shows that $|X_{1,}(\mathrm{t})| + |X_2(\mathrm{t})| \to 0$ as $\mathrm{t} \to \infty$ and the proof is complete.

As an example let us consider the case of two species Lotka- Volterra type competition in a periodic environment for for which we have

$$f_1\big(\mathrm{t}, x_1(\mathrm{t}), x_2(\mathrm{t} - T_{12})\big)1 \quad \left\{ \frac{x_1(\mathrm{t}) + ax_2(\mathrm{t} - T_{12})}{K_1(\mathrm{t})} \right\}$$

$$f_2(\mathrm{t}, x_1(\mathrm{t} - {}_{21})x_2(\mathrm{t})) = 1 - \left\{ \frac{\beta x_1(\mathrm{t}-T_{21}) + x_2(\mathrm{t})}{K_2(\mathrm{t})} \right\}$$

Where $\alpha, \beta$ are positive constants and $k_{1,}k_{2,}$ are strictly positive continuous periodic functions with a common period Equation (4.6)and (4.7) for this special case lead to

$$\left|\operatorname{axm}_{t>T}\{-a_{11}(t)\}\right| = \max_{t>T}(1/K_1(t)) > \min_{t>T}\left(-a_{21}(\mathrm{t})\right)\Big| = \min_{t\in}, \beta/K_2(\mathrm{t})$$

$$\max_{t\in T}\{-a_{22}(t)\}\Big| = \max_{t\in T}(1/K_2(t)) > \min_{t\in T}\left(-a_{12}(\mathrm{t})\right)\Big| = \min_{t\in} \mathfrak{a}/K_1(\mathrm{t})$$

The inequalities reduce in the absence of temporal variations in $k_{1,}k_{2,}$ to the inequality 0< $\mathfrak{a}$,$\beta$<1, which is well known for the global asymptotic stability of the positive steady state corresponding to the above example in a constant systemic domain the above example shows also that delays in interspecific interactions have no influence on the local asymptotic stability of the unique periodic solution of the periodic ecosystem A number of cases of this type in a constant environment have been established recently by the author (Gopalsamy [6]).A numerical example is illustrated at the end.

**Consequences**

It can be found that all the conclusions on the existence uniqueness and stability of a periodic solution of (2.1)-(2.4) are valid when $OT_{12}$=0 or $T_{12}$=0 or both $T_{12}$=0 and $T_{21}$=0: thus it will follow that delays in interspecific competition do not destabilize an otherwise stable coexistence (oscillatory or nonoscillatory) of two competing types of interacting particles atoms molecules transactions interaction. In forthcoming work we will extend the analysis of the system (2.1)-(2.4) to a system of n-competing species with a hypothesis similar to that of (2.4) It is however an open problem to investigate the effects of delays in intraspecific competition in a periodic milieu

Pianka ([8],pages 139-140) remarks that “In real population rate of increase competitive abilities and carrying capacities do very from individual to individual with population density in space and time Indeed temporal variation in the environment may often allow coexistence by continually altering the competitive abilities of population inhabiting it Moreover time lags are doubt lessly of some importance in real population Finally a heterogeneous environment may allow real competitors to evolve divergent resource utilization patterns and to reduce interspecific competitive inhibition

To understand the mechanism of selection pressure against competitive inhibition we first write the Malthusian parameters $\mu_1$, $\mu_2$ of fitness for the two species as follows (see MacArthur[7],page 226-230):

$$\mu_1=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_1\left(t,x_1(t),x_2(t-T_{12})\right)$$

$$\mu_2=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_2(t,x_1(t-T_{21}),x_2(t))$$

Since $\left(x_1(t),x_2(t)\right)\rightarrow\left(x_1^*(t),x_2^*(t)\right)$ for large t we can write the parameters of finess in the from

$$\mu_1=f_1\left(t,x_1^*(t),x_2^*(t-T_{12})\right)$$

$$\mu_2=f_2\left(t,x_1^*(t-T_{12}),x_2^*(t)\right)$$

In these relations we can use

$$x_2^*(t-T_{12}),x_2^*(t)-T_{12}\frac{dx_2^*(t-T_{12}+\theta_1T_{12}}{dt}$$

$$x_1^*(t-T_{12}),x_1^*(t)-T_{21}\frac{dx_1^*(t-T_{21}+\theta_2T_{21}}{dt}$$

$\theta_1,\theta_2,\in(0,1)$ and rewrite $\mu_1,\mu_2$ in the form

$$\mu_1=f_1\left(t,x_1^*(T),x_2^*(t)-T_{12}\frac{dx_2^*(t-T_{12}+\theta_1T_{12}}{dt}\right.$$

$$\mu_2=f_2\left(t,x_1^*(T),-T_{21}\frac{dx_1^*(t-T_{21}+\theta_2T_{21}}{dt},x_2^*(t)\right)$$

$\theta_1,\theta_2,\in(0,1)$

It is seen that the presence of the delays $T_{12}$ and $T_{21}$ provides additional sources of variability in the finess parameters and thereby in the selection pressure against interspecific competition Since a necessary and sufficient condition for natural selection to act is differential (variable) rates of confederation growth we find that time delays coupled with periodicity of an domain may allow real to evolve divergent resource utilization patterns and to reduce

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|---|
| Interspecific competitive inhibition for some remarks relating to the significance of a periodic environment to sustain the coexistence of a large number of confederation types in the absence of delays we refer to a recent article (Gopalsamy)The following example has been numerically solved and the solutions are displayed graphically for two different values of the delay parameter Τ: |
| $\frac{dx_1(\mathrm{t})}{dt} x_1(\mathrm{t})[1 - x_1(\mathrm{t}) + 0.4x_2(\mathrm{t} - \mathcal{T})/k_1(\mathrm{t})]$ |
| $\frac{dx_2(\mathrm{t})}{dt} x_2(\mathrm{t})[1 - \{0.8\, x_1(\mathrm{t} - \mathcal{T}) + x_2(\mathrm{t})\}/k_2(\mathrm{t})]$ |
| Where $k_1(\mathrm{t}) = 20 + \sin\ (\pi \mathrm{t})$ and $k_2(\mathrm{t})$=2.5+sin (πt) |
| @@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@@ |
| There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies |

systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational

content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification.

## GOPALSAMY'S MODEL APPIED TO COSMOLOGICAL PROBLEMS

## GLOBAL ASYMPTOTIC STABILITY IN AN ALMOST PERIODIC LOTKA- VOLTERRA SYSTEM

**Abstract**

**Gopalsamy studied** Sufficient conditions are obtained for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of a Lotka- Volterra system with almost periodic coefficients. Research was published Australian. Math Soc. B.27 (1986) 346-360. Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of

decay rate of perturbations. Sole intention, main objective and primary concern of the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are counted in the same vein as that of predator prey. Note that the model is applicable to each and every module. Quintessentially each parameter should have been given different name such as superscript 1,2, etc., as that has been done in the foregoing and a formal concatenation done at the end. Due to spatial restrictions we have not been able to do the line of thought and manner of analysis as has been done hereinbefore. This is the most important aspect of the model. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 7. Introduction

At the outset we assume that the system. This article is concerned with the derivation of a set of "easily verifiable" sufficient conditions for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of the Lotka- Volterra system

$$\frac{dx_1(t)}{dt} = x_i(t)\left\{b_1(t) - \sum_{j=1}^{n} a_{ij}\,(t)x_j(t)\right\};$$

$$t>t_0;\ t_0\ \in (-\infty,\infty);\ i = 1,2,\dots,n, \qquad (1.1)$$

where $b_1, a_{ij}\ (i,j = 1,2,\dots.,n)$ are nonnegative almost periodic functions defined for $t \in (-\infty,\infty)$.

In question satisfies Haimovici's equation based on age wise stratification given in the foregoing. Note that Haimovici's equations can be applied to Lotka-Volterra system term by term.

Let R and $R_n$ respectively denote the set of all real numbers and the $n$ - dimensional real Euclidean space; $R_n^+$ will denote the nonnegative cone of $R_n$. Since almost- periodic functions defined on R are bounded, we define the constants $b_i^l, b_l^u, a_{ij}^l, a_{ij}^u, (i,j = 1,2,\dots,n)$ by the following:

$$\inf_{t\in R} b_l(t) = b_i^l; \qquad \sup_{t\in R} b_l(t) = b_l^u, \qquad i,j = 1,2,\dots,n.$$

$$\inf_{t\in R} a_{lj}(t) = a_{lj}^l; \qquad \sup_{t\in R} a_{lj}(t) = a_{ij}^u.$$

We shall study the almost- periodic system (1.1) with the following assumption on the coefficients of (1.1);

$$b_l^l > 0; \qquad a_{ll}^l > 0, \tag{1.2}$$

$$b_l^l > \sum_{\substack{j=1 \\ j\neq i}}^{n} a_{lJ}^u (b_j^u / a_{iJ}^l). \tag{1.3}$$

Since the solutions of (1.1) corresponding to nonnegative initial conditions remain nonnegative subsequently, it will follow that

$$\frac{dx_1}{dt} \leq x_l\{b_l^u - a_{ll}^l x_i\}; \qquad t > t_0, i = 1,2,\dots,n, \tag{1.4}$$

As a consequence of which we will have

$$0 < x_l(t_0) \leq b_l^u / a_{ll}^l = x_i^u \Rightarrow x_i(t) \leq b_l^u / a_{Il}^l, \qquad i = 1,2,\dots,n; t \geq t_0. \tag{1.5}$$

(1.4) And (1.5) together lead to

$$\frac{dx_1}{dt} \leq x_l \left\{ b_l^u \sum_{\substack{j=1 \\ j\neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right) - a_{ll}^l x_i \right\}, \qquad i = 1,2,\dots,n; \tag{1.6}$$

Now (1.2), (1.3) and (1.6) lead to

$$x_i(t_0 \left( b_i^l - \sum_{\substack{j=1 \\ J\neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right) \right) (1/a_{ll}^u) = x_i^l$$

$$\Rightarrow x_l(t) \to x_i^l \ for\ t \to t_0,\ i = 1, 2, \dots, n. \quad (1.7)$$

8. It is our aim to show that the mode together with other models stated in the foregoing together with haimovicis model is applicable to whole lot of consummative debit – credit systems in cosmology, astrophysics, finance, economics, sociology, anatomy computation. We will not however divide the system studied here into 9 components/ modules due to restrictions on space….

From the foregoing preparation we have the following:

LEMMA 1.1. *Assume that* $b_i, a_{ij} (i, j = 1, 2, \dots, n)$ *are scalar nonnegative almost periodic functions defined on* ***R*** *such that (1.2) and (1.3) hold. Then the set S defined by*

$$S = \{x = (x_1, x_2, \dots, x_n) \in \boldsymbol{R}_n | x_* \to x_l \to x^*; i = 1, 2, \dots, n\} \quad (1.8)$$

*Where*

$$x_* = \min_{1 \to n} x_l^l; \qquad x_* = \min_{1 \to n} x_i^u \quad (1.9)$$

*is invariant with respect to (1.1)*

**Appellation, cognomen, compellation, designation, denomination of the conducive, infelicitous and unfortutious** responses to environment change may be influenced by changes in available domain and milieu under consideration. An example is the following study: Authors describe the gravitational degrees of freedom of the Schwarzschild black hole by one free variable. Authors introduce an equation which we suggest to be the Schrödinger equation of the Schwarzschild black hole corresponding to this model. We solve the Schrödinger equation explicitly and obtain the mass spectrum of the black hole as such as it can be observed by an observer very far away and at rest relative to the black hole. Our equation implies that there is no singularity inside the Schwarzschild black hole, and that the black hole has a certain ground state in which its mass is non-zero. **Schrodinger equation of the Schwarzschild black hole Jarmo Mäkelä (Cambridge U E-Print: gr-qc/9602008 | PDF** Particles, axions, atoms, molecules, quarks, Majorana fermions, bosons et al in its confederation capacity and aggregation processes, type interactions and interactions between landscape dynamics. Following bears ample testimony, infallible observatory to the fact there might be changes in the quantum behaviour due to merger of black holes. Following example bears enough testimony, infallible observatory and impeccable demonstration and apodictic evidence as how some of the models could be applied to cosmological problems. This is attributable and ascribable to the fact that the classification criterion is based on the age.

**How gravity kills (e) Schrödinger's cat Theorists argue that warped space-time prevents (e) quantum superpositions of large-scale objects. Elizabeth Gibney 17 June 2015 Propositions are written in model able form and can easy be scheduled.**

(67) If the cat in Erwin Schrödinger's famous thought-experiment behaved according to quantum theory, it would be able to exist in multiple states at once: both dead and alive.

(68) Physicists' common explanation for why we don't (e) see such quantum superpositions — in cats or any other aspect of the everyday world — is interference from (e) the environment.

(69) As soon as a quantum object interacts with (e&eb) a stray particle or a passing field, it picks (e) just one state, collapsing (eb) into our classical, everyday view.

(70) But even if physicists could completely isolate (eb) a large object in (eb) a quantum superposition, according to researchers at the University of Vienna, it would still collapse into (e&eb) one state — on Earth's surface, at least.

(71) "Somewhere in interstellar space it could be that the cat has (e) a chance to preserve quantum coherence, but on Earth, or near any planet, there's (=) little hope of that," says Igor Pikovski. The reason, he asserts, is gravity.

(72) Pikovski and his colleagues' idea, laid out in a paper published in Nature Physics on 15 June1, is at present only a mathematical argument. But experimenters hope to test whether gravity really does collapse (eb) quantum superpositions, says Hendrik Ulbricht, an experimental physicist at the University of Southampton, UK. "This is a cool, new idea, and I'm up for trying to see it in experiments," he says. Assembling the technology to do so, however, may take as long as a decade, he says.

(73) Cinema-goers who saw the film Interstellar are already familiar with the basic principle behind the Vienna team's work. Einstein's theory of general relativity states that an extremely massive object causes clocks near it to (e) run more slowly because (e) its strong gravitational field stretches (e&eb) the fabric of space-time (which is why a character in the film aged only an hour near a black hole, while (e) seven years passed on Earth). On a subtler scale, a molecule placed nearer the Earth's surface experiences (e&eb) a slightly slower clock than one placed slightly further away.

(74) Because of gravity's effect on (e&eb) space-time, Pikovski's team realised that variance in a molecule's position will also influence (e&eb) its internal energy — the vibrations of particles within (eb) the molecule, which evolve (eb) over time.

(75) If a molecule were put in a quantum superposition of two places, the correlation between position and (e&eb) internal energy would soon cause (eb) the duality to 'decohere' to (e) the molecule taking just one path, they suggest.

(76) "In most situations decoherence is due to (e) something external; here it's as though the internal jiggling is interacting with (e&eb) the motion of the molecule itself," adds Pikovski.

**A practical limit**

(77) No one has yet seen this effect because other sources of decoherence — such as magnetic fields, thermal radiation and vibrations — are (=) typically much stronger and cause (eb) quantum systems to collapse long before gravity becomes an issue. But experimenters are keen to try.

(78) Markus Arndt, an experimental physicist also at the University of Vienna, already tests whether quantum superpositions can be observed for (e) large, though not cat-sized, objects (see Fattening up Schrödinger's cat).

(79) He sends large molecules through (e&eb) a matter-wave interferometer, a system that gives (eb) each molecule a choice of (e) two different paths.

(80) In the classical view, a molecule travels down just one path; a quantum molecule effectively passes through (e&eb) both routes at once and interferes with (e&eb) itself to create (eb) a characteristic wave-like pattern (see image, below).

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|---|
| (81) A quantum interference pattern built up from complex molecules known as phthalocyanines.(deleted due to spatial constraints) |
| (82) A similar set up could be used to (e) test for gravity's ability to destroy (e) quantum behaviour: by comparing a vertical interferometer, in which the superposition should soon decohere due to (e) the stretching of time in one path versus (e&eb) the other, with a horizontal set up, where (e) the superposition could remain. |
| (83) Arndt, who has tested the effect for (e) molecules of up to 810 atoms2, points out those large molecules, would be (=) good for testing the gravitational effect because (e) they contain (e) lots of particles contributing to (e) internal energy. |
| (84) **But not only would the researchers have to further suppress the external environment to reduce (e) other decoherence effects, they would also need to (e) increase the separation of the two paths from micrometres to metres, or else use molecules perhaps 1 million times more massive. "It's certainly very challenging," says Arndt**. |
| (85) If gravity's effect does limit (e) quantum behaviour on Earth, tests of quantum reality for (e) large objects may eventually have to (e) move into space, says Angelo Bassi, a physicist at the University of Trieste in Italy. |
| (86) "But from a deep, fundamental point of view, this is nothing new," he says. A gravitational field is merely another environment to interact with, so invoking it does not (e) explain whether quantum behaviour might lead to (eb) classical reality if gravity's influence were mitigated — for example, by doing the experiment in (eb) gravity-free space. |
| (87) Nor does the effect described by Pikovski and colleagues say anything about quantum gravity: a theory that would unite gravity and (e&eb) quantum mechanics in (eb) a single description, which many researchers are working towards. |
| (88) "It's an interesting effect, but it's still quantum physics applied to (e&eb) classical general relativity. In that way it doesn't change our picture of the world," adds Bassi. Nature doi:10.1038/nature.2015.17773 |

Integrated factors these influences because they deal with ambience, mise en scene change or confederation dynamics, but rarely both. Time series of (e) space time suitability models with spatially explicit stochastic models to explore factors that influence the viability of type viability and confederation feasibility under stable and changing scenarios. Methodology could be applied to interrogation propositions and problems which might subsume one with the other notwithstanding that there is a space of singularities and there is space for completion of present from the past and future. An example is Time image movement from Deleuze. We clearly face a stiff challenge in our attempt to proffer Deleuzian treatments of photography as our critical project brushes up against the fact that Deleuze does not invest photography with (e&eb)the creative and re-creative potential that he associates with (e&eb) both cinema and modern painting. In short, while he highly values other visual artistic forms, he seemingly presents photographic texts as stagnate documents or tools that produce certainty, organize bodies and desires, and iterate hackneyed ideas. He even uses photography as something of a foil to demonstrate the innovation of the cinema and the originality of modern painters. In Francis Bacon: The Logic of Sensation (1981), he notes how "photography has taken over the illustrative and documentary role, so that modern painting no longer needs to fulfill this function, which still burdened earlier painters" (10). He later explains that "photographs are ways of seeing, and as such, they are illustrative and narrative reproductions or representations ... . They are what are seen, until finally one sees nothing else." He treats the photo as an instrument for reproducing representations of reality—a device that iterates images until they are ossified as established stories, icons, or even stagnant perceptions. Indeed, like Bergson before him, Deleuze compares photography to what he identifies as a reductionist model of perception: "what we see, what we perceive, are photographs" (74).[4] The photo is, in effect, an always-already passé sensational experience that numbs our sensitivity to the ongoing vitality of life and provides us with an efficient

understanding of the dynamism of our world. But it is compelling that Deleuze theorizes both the quotidian nature and the liberatory potential of photography; he treats it as a common, everyday visual experience that, despite its aesthetic and technological limitations, emancipates the modern painter to explore new artistic opportunities. This special issue of Rhizomes repeatedly provides artistic examples and critical strategies that invite us to see photography as replete with the power to unburden and emancipate that Deleuze associates with both film and painting. Indeed, one way to think of this project of reading Deleuze against Deleuze might be to ask: if photography can free painting and painters to pursue novel creative opportunities, can it likewise free itself, photographic practitioners, and even viewers to create and re-create anew? **Rhizomes » Issue 23 (2012) » Michael Kramp Unburdening Life, or the Deleuzian Potential of Photography Michael Kramp Lehigh University** dmk209@lehigh.edu

**Existence of an almost- periodic solution**

DEFINITION 2.1. A function $f(t,x)$,where $f$ is an $m$- vector, $t$ is a real scalar and $x$ is an $n$- vector, is said to be almost periodic in $t$ uniformly with respect to $x \in X \subset R_n$, $if\ f(t,x)$ is continuous in $t \in R\ and\ x \in X$, and if for any $\varepsilon > 0$, it is possible to find a constant $l(\varepsilon) > 0$ such that in any interval of length $l(\varepsilon)$ there exists a $\tau$ such that the inequality

$$\left|\left|f\ (t+\ \tau,x) - \ f(t,x)\right|\right| = \sum_{l=1}^{m} |f_l\ (t+\ \tau,x) = \ f_l(t,x)|$$

$$< \varepsilon \qquad (2.1)$$

is satisfied for all $t \in (-\infty,\infty), x \in X$. The number $\tau$ is called an $\varepsilon -$ translation number of $f(t,x)$.

Let $\{\lambda_j\}$ denote the set of all real numbers such that

$$\lim_{T\to\infty} T^{-1} \int_0^T f(t,x)\exp(-i\,\lambda_j\,t)dt \not\equiv 0 \quad for\ x \in X; i = \sqrt{-1}. \qquad (2.2)$$

It is known that when the set X is separable, the set of numbers $\{\lambda_j\}$ in (2.2) is countable. The set $\{\sum_1^N n_j\ \lambda_j\}$ for all integers N and integers $n_j$ is called the module of $f(t,x)$.

The following result concerned with module containment will be used in the proof of our existence theorem below.

**MODULE CONTAINMENT THEOREM** (Yoshizawa [7], page 18). *Let $f(t,x)$ and $g(t,x)$ be almost periodic in t uniformly for $x \in D \subset \boldsymbol{R}_n$. If for any compact set $S \subset D$ and for any sequence of real numbers $\{\tau_k\}$ having a limit as $k \to \infty$ (the limit being finite or$\pm\infty$) for which $\{f(t+\tau_k,x)\}$ is uniformly convergent on $\boldsymbol{R} \times$*

$S, \{g(t+,x)\}$ *is also uniformly convergent on* $\boldsymbol{R} \times S$ *then the module of* $g(t+,x)$ *is contained in the module of* $f(t+,x)$.

We can now formulate our existence theorem.

**THEORM 2.1** *Assume that the almost-periodic parameters* $b_1, a_{1J}$ $(i,j = 1,2,\dots,n)$ *of (1.1) satisfy (1.2)-(1.3). Furthermore assume that there exists a positive number* $\mu$ *such that*

$$\inf_{1 \in \boldsymbol{R}} a_{11}(t) > \sum_{\substack{J=1 \\ J \neq 1}}^{n} (\sup_{t \in \boldsymbol{R}} a_{Jl}(t)) + \mu, \quad i = 1,2,\dots,n. \tag{2.3}$$

*Then the system (1.1) has a strictly positive (component wise) almost-periodic solution, say* $w(t) = w_1(t), \dots, w_n(t), t \in \boldsymbol{R}$, *whose module is contained in that of*

$$f(t,x) = \{f_1(t,x_1,\dots,x_n), f_2(t,x_1,\dots,x_n), \dots, f_n(t,x_1,\dots,x_n)\}$$

*Where*

$$f_l(t,x_1,\dots,x_n) = x_l \left\{ b_l(t) - \sum_{j=1}^{n} a_{lj}(t) x_J \right\},$$

$$i = 1,2,\dots n; (x_1,x_2,\dots,x_n) \in S. \tag{2.4}$$

PROOF. The set $S$ defined by (1.8)-(1.9) is compact in $\boldsymbol{R}_n$. Let $\{\tau_h\}$ be any sequence of real numbers such that $\tau_k \to \infty$ as $k \to \infty$ and

$$\sum_{l=1}^{n} |f_i(t+\tau_k, x_1, \dots, x_n) - f_l(t, x_1, \dots, x_n)| \to 0 \quad \text{as } k \to \infty \tag{2.5}$$

Uniformly in $(x_1, \dots, x_n) \in S$ for $t \in (-\infty, \infty)$. for any real number $\beta$let $k_0(\beta)$ be the smallest value of $k$ such that $\tau k_0 + \beta \geq t_0$. Since the invariance of $S$ implies that any solution of (1.1)

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), x_2(t+\tau_k, t_0, x_2(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\} \in S$$

For all $t \geq \beta$ and $k \geq k_0(\beta)$ whenever $\{x_1(t_0), \dots, n_n(t_0)\} \in S$, our first task in the proof is to show that the sequence

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Converges to a continuous bounded function say $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined for $t \in [\beta, \infty)$ such that $w_1(t) \in [x_*, x^*], t \in [\beta, \infty), i = 1, 2, \dots, n$ and the convergence is uniform on all compact subsets of $[\beta, \infty)$. Since $x_*$ and $x^*$ are independence of $\beta$, it is enough to show that

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Is a Cauchy sequence on compact subsets of $[\beta, \infty)$.

Let $U$ be any compact subset of $\beta, \infty)$., and let be any arbitrary positive number. Choose an integer $n_0(\varepsilon, \beta) \geq k_0(\beta)$ so large that for $m \geq k \geq n_0(\varepsilon, \beta)$

and all $t \in (-\infty, \infty)$

$$\sum_{i=1}^{n} |b_1(t+\tau_m) - b_1(t+\tau_k)| < (\mu\varepsilon x_*)/(4x^*), \tag{2.6}$$

$$x^* \sum_{i=1}^{n}\sum_{j=1}^{n} |a_{lj}(t+\tau_m) - a_{lj}(t+\tau_k)| < (\mu\varepsilon x_*)/(4x^*) \tag{2.7}$$

$$(2nx^*/x_*)\exp[-\mu x_*(\beta + \tau_k - t_0)] < (\varepsilon/2x^*)$$

(2.8)

Denote $x_i(t, t_0, x_1, (t_0))$ $by$ $x_1(t)$ brevity, and consider a function $v(s)$ defined by

$$v(s) = v\left(x(s), x(s+\tau_m+\tau_k)\right)$$

$$= \sum_{i=1}^{n} |\log x_1, (s) - \log x_1(s+\tau_m+\tau_k)|, \qquad x \geq t_0, \tag{2.9}$$

Where $x(t) = \{x_1(t), \dots, x_n(t)\}$ denotes a solution of (1.1) with $\{x_1, (t_0), \dots, x_n(t)\} \in S$. the invariance of S for (1.1), together with the elementary mean value theorem of differential calculus, implies

$$(1/x^*)\sum_{l=1}^{n}|x_l(s)-x_1(s+\tau_m+\tau_k)|\ \upsilon(s)$$

$$(1/x^*)\sum_{i=1}^{n}|x_l(s)-x_1(s+\tau_m+\tau_k)|. \tag{2.10}$$

Calculating the right derivative $D^+\upsilon$ of $\upsilon$ along the solutions of (1.1) and simplifying

$$D^+\upsilon(s)=\sum_{i=1}^{n}\{sign[x_i(s)-x_i(s+\tau_m-\tau_k)]\}$$

$$\left\{\frac{d[logx_i(s)]}{ds}-\frac{d[logx_i(s+\tau_m-\tau_k)]}{ds}\right\}$$

$$\sum_{l=1}^{n}|b_l(s)-b_l(s+\tau_m-\tau_k)|$$

$$+\sum_{l=1}^{n}(-a_{il}(s)|x_i(s)-x_i(s+\tau_m-\tau_k)|+\sum_{\substack{J=1\\ j\neq l}}^{n}a_{ij}(s)|x_j(s)-x_j(s+\tau_m-\tau_k)|)$$

$$+\sum_{i=1}^{n}\sum_{j=1}^{n}|a_{ij}(s)|a_{ij}(s+\tau_m-\tau_k)|\,x_j(s+\tau_m-\tau_k) \tag{2.11}$$

$$-\mu x_*\upsilon\big(x(s),x(s+\tau_m-\tau_k)\big)+\sum_{l=1}^{n}|b_l(s)-b_i(s+\tau_m-\tau_k)|$$

$$+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|. \tag{2.12}$$

Let us set $s-\tau_k=\sigma$; then we have from (2.6)-(2.7) that

$$\sum_{i=1}^{n}|b_l(s)-b_1(s+\tau_m-\tau_k)|+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|.$$

$$= \sum_{l=1}^{n} |b_l(\sigma + \tau_k) - b_i(\sigma + \tau_m)|$$

$$+x^* \sum_{i=1}^{n} \sum_{J=1}^{n} |a_{ij}(\sigma + \tau_k) - a_{ij}(\sigma + \tau_m)|$$

$$< (\mu\varepsilon x^*)/(2x^*). \quad (2.13)$$

Now an integration of (2.12) over $(t_0, t - \tau_k)$ together with (2.13) leads to

$$v\big(x(t+\tau_k), x(t+\tau_m)\big) \leq v\big(x(t_0), x(t_0 + \tau_m - \tau_k)\big) \exp[-\mu x_*(t + t_k - t_0)]$$

$$+[(\mu\varepsilon x_*)/(2x^*)]/\mu x_*)$$

$$\leq [(2x_n^*)/x_*]\exp[-\mu x_*(t + t_k - t_0)] + (\varepsilon/2x^*)$$

$$< (\varepsilon/x^*) \text{ On using (2.8), (2.10)} \quad (2.14)$$

For $m \geq n_0(\varepsilon, \beta)$ and for all $t \in \cup \subset [\beta, \infty]$, Thus we have from the definition of $v$ and (2.14) that

$$\sum_{l=1}^{n} |log x_i(t + \tau_k) - log x_l(t + \tau_m)| < (\varepsilon/x^*) \quad (2.15)$$

And hence using $x_* \leq x_i(t) \leq x^*, i = 1, 2, \dots, n; t \geq t_0$,

$$\sum_{i=1}^{n} |x_i(t + \tau_k) - x_i(t + \tau_m)| < \varepsilon \text{ for all } t \in \cup \subset [\beta, \infty], m \geq n_0(\varepsilon, \beta). \quad (2.16)$$

The existence of a function $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined on $[\beta, \infty]$, such that $w_i(t) \in [x_*, x^*], i = 1, 2, \dots, n; t \in [\beta, \infty]$ follows from (2.16) and the fact that $x_i(t) \in [x_*, x^*], i = 1, 2, \dots, n;$ for $t \geq t_0$. Since $\beta$ is arbitrary, $w$ is in fact

Defined on $(-\infty, \infty)$ and hence we have

$x_i(t+\tau_k) \to w_l(t)$ as $k \to \infty$ uniformly on compact subsets of $(-\infty,\infty)$, $i = 1,2,\dots,n.$

(2.17)

Our next task is to show that $w: R \to S$ is differentiable and is a solution of (1.1). Since $x(t) = \{x_1(t),\dots,x_n(t)\}$ satisfies (1.1) we have $\sum_{i=1}^{n} |\frac{d}{dt} x_i(t+\tau_k) - \frac{d}{dt} x_i(t+\tau_m)|$

$$\sum_{i=1}^{n} |f_l\left(t+\tau_k, x_1(t+\tau_k),\dots, x_n\ (t+\tau_k)\right)$$

$$- f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)| \qquad (2.18)$$

$$+\sum_{l=1}^{n} |f_l\ (t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\ )$$

$$- f_l\left(t+\tau_m, x_1(t+\tau_m),\dots,x_n(t+\tau_m)\right)|$$

Where $f_l\ (i = 1,2,\dots,n)$are as in (2.4). Since $x(t+\tau_k) = x_1(t+\tau_k),\dots,\ x_n(t+\tau_k) \in S$ for large $\tau_k$ for $t$ in each compact subset of $(-\infty,\infty)$ there exists an integer $n_1(\varepsilon) > 0$ such that if $m$ ? $k$ ? $n_1(\varepsilon)$, then for any $\varepsilon > 0$,

$$\sum_{i=1}^{n} |f_l\left(t+\tau_k, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$-f_1\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)| < (\varepsilon/2). \qquad (2.19)$$

Similarly since $x_l(t+\tau_k) \in S$ and $x_l(t+\tau_m) \in S$ for $i = 1,2,\dots,n$ there is an integer $n_2(\varepsilon) > 0$ such that for large enough $m, k$, such that $m$ ? $k$ ? $n_2(\varepsilon)$, we have

$$\sum_{l=1}^{n} |f_l\left(t+\tau_m, x_1(t+\ \tau_k),\dots,x_n(t+\ \tau_k)\right)$$

$$-f_1\left(t+\tau_m, x(t+\tau_m),\dots,x_n(t+\tau_m)\right)| < (\varepsilon/2) \qquad (2.20)$$

This equation will be very much useful in chemistry if used with the modules mentioned and written in concatenation with each other. It should be borne in mind that the applicability is modulewise and in no case concatenation should be marginalized. In all cases parameters are different which could be represented with different superscripts and subscripts like 1,23,…….etc., **Anal Bioanal Chem. 2004 Sep; 380(1):36-45. Epub 2004 Jul 29. Electronic tongue for pharmaceutical analytics: quantification of tastes and masking effects Legin A et al**

The organoleptic aspects of pharmaceutical formulations affect their acceptability to the patient and hence can have an important effect on concordance with treatment. Objective evaluation of these aspects, particularly the taste of the formulation and the drug substance it contains, is difficult. Whilst volunteer taste panels can be used to good effect their utility is limited, particularly during very early stage development when the toxicological profile of the active pharmaceutical ingredient (API) is yet to be established in detail. A potentiometric "electronic tongue" has been applied to analyse a variety of 41 individual substances and mixtures of particular interest for pharmaceutical research and development. The electronic tongue (ET) was capable of discriminating between substances with different taste modalities and could also distinguish different substances eliciting the same basic taste; the ET is promising in terms of quantifying the content of each substance and has an ability to detect nuances of the basic taste (e.g. lingering or short-lived). After calibration the electronic tongue was successfully applied to predicting bitterness strength of binary mixtures with a sweetener in terms of "apparent" or "perceived" quinine content. In order to render a formulation palatable it is often necessary to mask the (usually bitter) taste of the API by the addition of masking agents such as sweeteners and flavours. The ET proved capable of distinguishing between formulations with different levels of sweetener and/or flavour in a manner that was consistent with their masking efficiency as perceived by a small human taste panel. A suitably calibrated ET could have the benefit of providing the pharmaceutical formulator with reliable data concerning the taste of the product quickly and with a reduced need to ask volunteers to taste active pharmaceutical samples. Early development activities could be facilitated when human tasting is usually not possible in the absence of the required toxicological data.

Thus if $m \neq n_0(\varepsilon) = max\{n_1, n_2\}$, we have

$$\sum_{l=1}^{n} |\frac{d}{dt}\, x_1(t+\ \tau_k) - \frac{d}{dt} x_1\ (\ (t+\tau_m)| < \varepsilon \qquad (2.21)$$

Which shows that

$$\lim_{k\to\infty} \frac{d}{dt} x_1(t+\ \tau_k) \quad exists\ for\ i = 1, 2, \dots, n$$

Uniformly on all compact subsets of $(-\infty, \infty)$. Thus we have

$$\lim_{k\to\infty} \lim_{k\to\infty} \frac{x_i(t+\tau_k+h) - \ x_i(t+\tau_k)}{h}$$

$$= \lim_{h\to 0} \lim_{k\to 0} \frac{x_l(t+\tau_k+h) - \ x_l(t+\tau_k)}{h} \qquad (2.22)$$

$$= \lim_{h\to 0} \frac{w_1(t+h) - \ w_1(t)}{h}, \qquad i = 1, 2, \dots, n,$$

Showing that (since the left side of (2.22) exists) $w_i (i = 1, 2, \dots, n)$ is differentiable. Furthermore we have from (2.22) that

$$\frac{dw_l(t)}{dt} = \lim_{k\to\infty} [\, f_i\ \big(t+\tau_k, x_1(t+\tau_k), \dots, x_n\ (t+\tau_k)\big)$$

$$-f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big)$$

$$+ - f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big)]$$

$$= \ -f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big); \qquad i = 1, 2, \dots, n, \qquad (2.23)$$

By the uniform almost periodicity of $f_1, f_2, \dots, f_n$ on $R \times S$ and hence $w(t) = \big(w_1(t), \dots, w_n(t)\big)$is a solution of (1.1).

In order to complete the proof, we have to show that $w(t) = \big(w_1(t), \dots, w_n(t)\big)$ is an almost- periodic function whose module is contained in that of $\{f_1(t, x_1, \dots, x_n), f_2(t, x_1, \dots, x_n), \dots, f_n(t, x_1, \dots, x_n)\}$ for $(x_1, x_2 \dots, x_n) \in S$. For this purpose it is in fact sufficient (by the module containment theorem) to show that for any sequence of real numbers $\{\tau_k\}$ for which

$$\{f_1(t + \tau_k, x_1, \dots, x_n), f_2(t, + \tau_k, x_1, \dots, x_n), \dots, f_n(t, + \tau_k, x_1, \dots, x_n)\}$$

Converges uniformly on R for all $(x_1, \dots, x_n) \in S$, the sequence $\{w_1(t, + \tau_k) \dots, w_n(t + \tau_k)\}$ converges uniformly on R where $\tau_k$ tends to a finite number or $\pm\infty$ as $k \to \infty$.

In the case where $\tau_k \to \tau \neq \pm\infty$ as $k \to \infty$, we have $w_l(t, + \tau_k) \to w_i(t + \tau), i = 1, 2, \dots, n$ uniformly on R as $k \to \infty$. Let us suppose $\tau_k \to \infty$ as $k \to \infty$. For any $\varepsilon > 0$ there exists an integer $n_0(\varepsilon)$ such that if $m$ [illegible] $n_0(\varepsilon)$ then

$$\sum_{l=1}^{n} |b_l(t + \tau_k) - b_l(t + \tau_m)| < (\varepsilon/4x^*)x_*, \qquad (2.24)$$

$$x^* \sum_{i=1}^{n} \sum_{j=1}^{n} |a_{ij}(t + \tau_k) - a_{ij}(t + \tau_m)| < \left(\frac{\varepsilon}{4x^*}\right)\mu x_* \qquad (2.25)$$

$$n(2x^*/x_*) \exp[\mu x_* \tau_k] < (\varepsilon/2x^*). \qquad (2.26)$$

Consider the function $v\big(w(s), w(s + \tau_m - \tau_k)\big)$ for $s \in [t, t + \tau_k]$ defined by

$$v\big(w(s), w(s + \tau_m - \tau_k)\big) = \sum_{l=1}^{n} |\log w_l(s) - \log w_l(s + \tau_m - \tau_k)|. \qquad (2.27)$$

Using the fact $w: \boldsymbol{R} \to S$ and $w$ is a solution of (1.1), calculate the right derivative $D^+v \ of \ v$;

$$D^+v\big(w(s), w(s + \tau_m - \tau_k)\big) = \sum_{i=1}^{n} \{\text{sign}[w_l(s) - w_i(s + \tau_m - \tau_k)]\}.$$

$$\left\{\frac{dw_i(s)}{ds} \cdot \frac{1}{w_l(s)} - \frac{dw_l(s + \tau_m - \tau_k)}{ds} \frac{1}{w_l(s + \tau_m - \tau_k)}\right\}$$

$$\leq -\mu \sum_{l=1}^{n} |w_l(s) - w_l(s+\tau_m-\tau_k)| + \sum_{l=1}^{n} |b_i(s) - b_l(s+\tau_m-\tau_k)|$$

$$+x^* \sum_{l=1}^{n}\sum_{J=1}^{n} |a_{lJ}(s) - a_{lJ}(s+\tau_m-\tau_k)| \quad (2.28)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+ \sum_{i=1}^{n} |b_l(s) - b_l(s+\tau_m-\tau_k)| \quad (2.29)$$

$$+x^* \sum_{l=1}^{n}\sum_{j=1}^{n} |a_{lJ}(s) - a_{ij}(s+\tau_m-\tau_k)|.$$

If we set $s-\tau_k = \sigma$ then we have from (2.29),

$$D^+ v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+ \sum_{l=1}^{n} |b_l(\sigma+\tau_k) - b_l(s+\tau_m)|$$

$$+x^* \sum_{l=1}^{n}\sum_{j=1}^{n} |a_{ij}(\sigma+\tau_k) - a_{ij}(\sigma+\tau_m)| \quad (2.30)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+[(\varepsilon\mu x_*)/(4x^*)] + [(\varepsilon\mu x_*)/(4x^*)]$$

(2.31)

Integrating (2.31) over $(t, t+T_k)$

$$v(w(t+T_k)w(t+T_m))$$

Another MODEL where the model is directly applicable with the scheduled ones are the case of insurance and corporate finance. Insurance: Mathematics and Economics Volume 42, Issue 1, February 2008, Retrieval of Black–Scholes and generalized Erlang models by perturbed observations at a fixed time Dedicated to Professor André Dubey on the occasion of his 60th birthday Daniel Neuenschwander

S-stable laws on the real line (more generally on Hilbert spaces), associated with some non-linear transformations (so-called "shrinking operations"), were introduced in [Jurek, Z.J., 1977. Limit distributions for truncated random variables. In: Proc. 2nd Vilnius Conference on Probability and Statistics, June 28–July 3, 1977. In: Abstracts of Communications, vol. 3, pp. 95–96; Jurek, Z.J., 1979. Properties of s-stable distribution functions. Bull. Acad. Polon. Sci. Sér. Math. XXVII (1), 135–141; Jurek, Z.J., 1981. Limit distributions for sums of shrunken random variables. Dissertationes Math. vol. CLXXXV]. In [Jurek, Z.J., Neuenschwander, D., 1999. S-stable laws in insurance and finance and generalization to nilpotent Lie groups. J. Theoret. Probab. 12 (4), 1089–1107], the authors interpreted s-stable motions on the real line as limits of total amount of claims processes (up to a deterministic premium) of a portfolio of excess-of-loss reinsurance contracts and showed that they led to Erlang's model or to Brownian motion. In [Neuenschwander, D., 2000b. On option pricing in models driven by iterated integrals of Brownian motion. In: Mitt. SAV 2000, Heft 1, pp. 35–39], we considered stochastic integrals whose integrand and integrator are both independent Brownian motions, thus modelling a stochastic volatility; as a result we got an analogue of the Black–Scholes formula in this model, confirming a result of Hull and White [Hull, J., White, A., 1987. The pricing of options on assets with stochastic volatility. J. Finance XLII (2), 281–300]. In the present paper, we will look at a common generalization of these processes, namely s-stable motions on the real line perturbed by a stochastic integral whose integrand and integrator are both (not necessarily independent) s-stable motions. The main result will be that if we can observe the distribution of such so-perturbed s-stable motions (together with the values of the perturbing processes) at time t=1, then we can identify the whole model (including the perturbation) among all models with Lévy processes perturbed by an iterated stochastic integral of two Lévy processes (in the Gaussian case) resp. among all models with a compound Poisson process with drift perturbed by an iterated stochastic integral of two compound Poisson processes (in the completely non-Gaussian case if the perturbing processes have no drift) without knowing anything about the history or about its distribution during $0 \le t < 1$. This applies, e.g., to a situation

where several assets obey the same model and one can estimate the distribution at time one by looking at the values of all these assets at time t=1. Interestingly enough, it will be convenient to treat the whole matter in the algebraic framework of the so-called Heisenberg group. This is a concept coming in fact from quantum mechanics and is in a certain sense the simplest non-commutative Lie group.

$\leq e^{-\mu x^* T_k} v\big(w(t), w(t+T_m-T_k)\big) + (\varepsilon/2x^*)$

$\leq e^{-\mu x^* T_k} \sum_{i=1}^{n} |\log\ w_i(t) - \log\ w_i(t+T_m-T_k)| + (1/x_*) + (\varepsilon/2x^*)$

$\leq n(2x^*/x_*)\exp[-\mu x_* T_k] + (\varepsilon/2x^*)$

$<(\varepsilon/x^*)$.

Since

$$v\big(w(t+T_k)w(t+T_m)\big) \leq \sum_{i=1}^{n} |\ w_1(t+T_k) - w_1(t+T_m))/x^*)$$

We have from(2.33)that

$$\sum_{i=1}^{n} |\ w_1(t+T_k) - w_1(t+T_m)| < \varepsilon\ for\ m \geq k \geq n_0(\varepsilon)$$

Which show that the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\}$converges uniformly on **R**. By a similar argument the above conclusion is also valid when $T_k \to -\infty\ as\ k \to \infty$. Thus the uniform convergence of the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\}\ as\ k \to \infty\ on\ \mathbf{R}$ follows and this completes the proof of the theorem.

**COROLLARY.** Suppose $b_i a_{ij}$ (i, j = 1.2, … , n)in (1.1)are periodic with a common period say $T > 0$. *Then if* (1.2), (1.3)*and* (2.3)*hold the system* (1.1)*will have a periodic solution* $w(t) = \{w,(t), \ldots\ldots, w_n(t)\}$*of period* $T$ *such that*

$$w_1(t+T) = w_1(t), \qquad t \in \boldsymbol{R}: i = 1,2, \ldots, n.$$

**PROOF.** Proof follows from that of Theorem 2.1 if we Choose $T_k = KT(K = 0,1,2, \ldots)$*and note that*

$$x_1(t+KT) - w_1(t) \to 0\ as\ k \to \infty; i = 1,2, \ldots, n; t \in \boldsymbol{R},$$

$$x_i(t+(K+1)t) - w_1(t) \to 0 \ \ as \ \boldsymbol{K} \to \infty; i = 1,2,\dots,n; t \in \boldsymbol{R},$$

$$x_i(t+T+KT) - w_1(t+T) \to 0 \ \ as \ k \to \infty; i = 1,2,\dots,n; t \in \boldsymbol{R},$$

Implying that

$$w_1(t+T) = w_i(t); \quad i = 1,2,\dots,n; t \in \boldsymbol{R},$$

## 3. Global asymptotic stability

By definition we say that an almost-periodic solution say $w(t) = \{w_1(t),\dots,w_n(t)\} of (1.1) is \ globally \ asymptotically$ stable (or attractive) if and only if every other solution x(t)={ $x_1(t),\dots,x_n(t)$}with $x_1(o) > 0, i = 1.2,\dots,n$ is defined for all t□0 and satisfies

$$\lim_{t\to\infty} |w_1(t) - x_1(t)| = 0; \quad i = 1,2,\dots,n.$$

A consequence of such a global asymptotic stability is that there cannot be another strictly positive almost-periodic solution of (1.1)

**THEOREM 3.1** Suppose the almost- Periodic coefficients $b_1$, $a_{ij}$ (i, j = 1,2, … , n)

Satisfy (1.2)and (1.3)and the following

$$\inf_{t\in R} a_{ij}(t) > \left\{ \sum_{\substack{J=1 \\ J\neq 1}}^{n} \sup a_{ij}(t) \right\} + \mu; \qquad i = 1,2,\dots,n$$

For some positive number $\mu \ Then$ (1.1)has an almost-periodic solution (with strictly positive components) which is globally asymptotically stable.

We will omit the proof since the proof is identical with that of Theorem 3.1 of [3]

Under condition (3.2) the almost –periodic solution $w(t) = \{w_1(t), \ldots . w_n(t)\}$ of (1.1)enjoys a somewhat strong from of stability even if the coefficients of (1.1)are subjected to a class of perturbations We note (see Yoshizawa [7]page17)that if a function b:R→ $R$ is almost periodic then H(b),the "hull of b" is defined as follows:

H(b)= g|g:**R**→ $\boldsymbol{R}$ $and\ g(t)\lim_{k\to\infty} b(t+T_k) for$ some sequence

$T_k$ of real numbers the limit being uniform in $t \in R$.

By Theorem 2.2 of Yoshizawa [7, page 10], members of the hull of almost periodic functions are themselves almost periodic.

A system of the from

$$\frac{dy_1(t)}{dt} = y_1(t)\left\{B_i(t) - \sum_{J-1}^{n} A_{ij}\ (t)y_1(t)\right\},$$

*i=1,2,…,n; t>*$t_0 \in R$,

Where

$$B_i \in H(b_i); (A_{ij}), \qquad \text{i, j} = 1,2, \ldots, \text{n},$$

is known as a system from the hull of (1.1) As a consequence of the definition of the hull in (3.3)and (3.5)we shall have

$$\inf_{t\in R} B_i(t) = \text{b}_i^l; \quad \sup_{t\in R} B_i(t) = \text{b}_i^l;$$

$$\inf_{t\in R} A_{ij}(t) = \text{b}_i^l; \quad \sup_{t\in R} A_{ij}(t) = \text{a}_{iJ}^u;$$

and hence the coefficients of(3.4)- (3.5)satisfy the same conditions as (1.2),(1.3)and (3.2)It will then follow from theorems(2.1)and (3.1)that solutions of (3.4)-(3.5)with $y_1(t_0) > 0, i = 1,2, \ldots, n\ will\ have$ the following property:

$$y_1(t) > 0 \ \ for\ t \geq t_0; i = 1,2,\dots,n,$$

$$\sum_{i=1}^{n} y_1(t) remains\ bounded\ for\ t \geq t_0$$

**DEFINITION** A Solution $x(t) = \{x_1(t), \dots, x_n(t)\}, t \in t_0 \in \boldsymbol{R}$,of (1.1)is said to be stable under perturbations from the hull of (1.1)if the following holds :let $y(t) = \{(y_1(t), \dots, y_n(t)\}$ be any solution of (3.4)-(3.5)such that $y_1(t_0) > 0$, i=1,2,…,n,and for some positive constants $\beta_1, \beta_2 \beta_1, \leq y_1(t) \leq \beta_2, t \geq t_0, \ \ i = 1,2,\dots,n$ Then given any $\varepsilon > 0$,there exists a constant $\delta > 0$ such that

$$\sum_{j=1}^{n} |x_1(t_0) - y_1(t_0)| < \delta$$

$$\sum_{i=1}^{n} \sup_{t \in R} |b_1(1) - B_1(t)| < \delta \quad = \sum_{i=1}^{n} |x_1(t) - y_i(t)| < \varepsilon$$

$$\sum_{i=1}^{n} \sum_{j=1}^{n} | \quad \sup_{t \in R} |a_{ij}\ (t) - A_{ij}(t)| < \delta$$

For $t \geq t_0$.

We can now formulate the following result which shows that all solution of (1.1)are stable under perturbations from the hull of (1.1)

**THEOREM 3.2**. Assume that the conditions of Theorem (3.1) hold. Then every strictly positive solution of (1.1) ( including its unique almost-periodic solution) is stable under perturbations from the hull of (1.1).

PROOF. Let $x(t) = \{x_1(t), \dots, x_n(t)\}$ and $y(t) = \{y_1(t), \dots, y_n(t)\}$ be any two strictly positive solutions of (1.1) and (3.4) – (3.5) respectively for $t \geq t_0 \in R$. Let $\beta_1, \beta_2$ be constants such that

$\beta_1 \leq x_i(t) \leq \beta_2$ ,
$\beta_1 \leq y_i(t) \leq \beta_2$ , $i = 1,2,\dots,n; t \geq t_0$ (3.7)

Consider now a function $V(t) = V(x(t)y(t))$ defined by

$$V\big(x(t)y(t)\big) = \sum_{i=1}^{n} \big|\log x_1(t) - \log, y_i(t)\big|, \quad t \geq t_0 \in \text{R}. \tag{3.8}$$

It is easy to see that

$$\frac{1}{\beta_2} \sum_{i=1}^{n} \big|x_i(t) - y_i(t)\big| \leq V\big(x(t)y(t)\big)$$

$$\leq \frac{1}{\beta_1} \sum_{i=1}^{n} \big|x_i(t) - y_i(t)\big|, \; t \geq t_0 \in \text{R}. \tag{3.9}$$

Calculating the right derivative $D^+V$ of $V$ and using (1.1) and (3.4) we derive (after simplification) that

$$D^+V\ \big(x(t), y(t)\big) \leq -\mu \sum_{i=1}^{n} \big|x_i(t) - y_i(t)\big| + \sum_{i=1}^{n} \big|b_i(t) - B_i(t)\big|$$

$$+\beta_2 \sum_{j=1}^{n} \sum_{i=1}^{n} \big|a_{ij}(t) - A_{ij}(t)\big|; \quad t \geq t_0,$$

$$\leq -\mu\, \beta_1\, V\big(x(t), y(t)\big) + \sum_{i=1}^{n} \big|b_i(t) - B_i(t)\big|$$

$$+ \beta_2 \sum_{i=1}^{n} \sum_{j=1}^{n} \big|a_{ij}(t) - A_{ij}(t)\big|; \quad t \geq t_0, \tag{3.10}$$

An integration of (3.9) over $[t_0, t]$ with an application of (3.8) leads to

$$\frac{1}{\beta_2}\sum_{i=1}^{n}|x_i(t)-y_i(t)| \leq V(x(t),y(t)) \leq V(x(t_0),y(t_0)) + \left[\sum_{i=1}^{n}\sup_{t\in R}|b_i(t)-B_i(t)| + \beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}|a_{ij}(t)-A_{ij}(t)|\right]\frac{1}{\mu\beta_1},$$

$$t \geq t_0,$$

And hence

$$\sum_{i=1}^{n}|x_1(t)-y_1(t)| \leq \left(\frac{\beta_2}{\beta_1}\right)\sum_{i=1}^{n}|x_1(t_0)-y_1(t_0)| + \frac{\beta_2}{\mu\beta_2}\left[\sum_{i=1}^{n}\sup_{t\in R}|b_i(t)-B_i(t)|\right.$$

The construction reason is because is criticized in a number of usage guides: The reason for the long delays was because the costs greatly exceeded the original estimates. One objection to this construction is based on its redundancy: the word because (literally, by cause) contains within it the meaning of reason; thus saying the reason is because is like saying "The cause is by cause," which would never be said. A second objection is based on the claim that because can introduce only adverbial clauses and that reason is requires completion by a noun clause. Critics would substitute that for because in the offending construction: The reason for the long delays in completing the project was that the costs.

$$\left.+\beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}|a_{ij}(t)-A_{ij}(t)|\right]. \qquad (3.11)$$

Now for any $\varepsilon > 0$ if one chooses a $\delta > 0$ such that

$$\delta < \min\left\{\frac{\varepsilon}{3}\frac{\beta_1}{\beta_2}, \frac{\varepsilon}{3}\frac{\beta_1}{\beta^2}\mu\right\} \qquad (3.12)$$

Then (3..6) will be satisfied and this completes the proof

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# . Jia Li's Prescientiial Dynamics And Contential Staticity Of Accentuation Attrition Models Applied To Cosmological And Astrophysical And Quantum Mechanical Engineering Problems

**Abstract**: **Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics ET AL ). We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems. Predator prey model literally form bastion, pillar, post, stylobate, sentinel towards the end of dismasting many of the asymptotic and solutional behaviour in finance, economic, and as we shall see in cosmology, quantum computation and other fields.**

## MODEL FORMULATION

Let $p_1(t, a_1)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_1}{dt} = -b_1 Q_1 + g_1(g_1)Q_1$$

Where $b$ is the natural death rate in the absence of credit types and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_1(0) = 0, \quad g_1(P_1) > 0, \; P_1 \; 0, \qquad \lim_{P_1 \to \infty} g_1(P_1) = (K_1)_1 \; \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_1 + (\rho_a)_1 + D_1(a_1, P_1, Q_1)\rho_1 + 0, \qquad 0 < a_1 \; \infty,$$

$$\rho_1(t, 0) = \int_0^{\infty} n_1(a_1, P_1, Q_1)\rho_1(t_1, a_1)da_1$$

Where $D_1$ is the mortality function and $n_1$ is the fertility function.

Put $_1$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_1}{dt}\right]_1 = x_1\left[\frac{dQ_2}{dt}\right]_1 - x_2\left[\frac{dQ_1}{dt}\right]_1$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_1 - x_4\left[\frac{dQ_2}{dt}\right]_1$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_1 - x_6\left[\frac{dQ_3}{dt}\right]_1$$

Where $x_1, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_1$ and $g_1(P_1)Q_1$

Now when you add the term $(b_1Q_1)$ with $g_1(P_1)(Q_1)$ we are infect see is that $(b_1Q_1)$ is credited incrementally part of $g_1(P_1)(Q_1)$ or instant nearly, when subtracted from $\frac{dQ_1}{dt}$, it means that $g_1(P_1)(Q_1)$ is slowly eaten by $\frac{dQ_1}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_1 = \mu_1(a_1) + (f_1)_1(P_1)_1 + (f_2)_1(Q)_1 \quad (2)$$

$$n_1 = \beta_1(a_1)h_1(Q_1), \quad (3)$$

Where $\mu_1$the age- dependent debit rate whit is $\mu_1(a_1)$ 0 $\forall a > 0$; $(f_1)_1$ is the density- dependent death rate with $f_1(0) = 0$, $(f_1^1)_1(P_1)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_1)_1(P_1) = (k_1)_1$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_1(0) = 0, (f_2^1)_1(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_1(Q_1) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^1(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_1 + (\rho_a)_1 = -\mu_1(a_1) + (f_1)_1P_1) + \big(f_2)_1(Q_1)\big)(\rho_1)_1 \quad (4)$$

$$\rho_1(t_1, 0) = h_1(Q_1)\int_0^\infty \beta_1(a_1)\rho_1(t, a_1)da, \quad (5)$$

$$\frac{dQ_1}{dt}$$

$$= -b_1Q_1 + g_1(P_1)Q_1 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_1$ and $Q_1$ of (4-6) satisfy

$$0 \leq \rho_1(t, a_1) \leq \rho_1^*(t, a_1),$$

$$0 \leq Q_1(t) \leq Q_1^*(t),$$

Where $\rho_1^*(t, a_1)$ and $Q_1^*(t)$, are solutions of

$$(\rho_t^*)_1 + (\rho_a^*)_1 = -(\mu_1(a_1) + (f_1)_1(P_1^*) + (f_2)_1 Q_1^*))\, \rho_1^*, \quad (7)$$

$$\rho_1^*(t, 0) = \int_0^\infty \beta_1(a_1)\, \rho_1^*(t, a_1) da_1 \quad (8)$$

$$\frac{dQ_1^*}{dt} = -b_1 Q_1^* + g_1(P_1^*) Q_1^*, \quad (9)$$

With $P_1^*(t) = \int_0^\infty \rho_1^*(t, a_1) da_1$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_1^*(t)}{dt} = P_1^*(t)\big(r_1 - (f_1)(P_1^*) - (f_2)_1(Q^*)\big), \quad (10)$$

$$\frac{dQ_1^*(t)}{dt} = Q_1^*(t)(-b_1 + g_1(P_1^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_1(a) e^{-\int_0^a \mu_1(\tau) dt} e^{-ra}\, da_1 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_1^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(t)dt}\, da_1 > 1, \qquad (H1)$$

Which guarantees $r_1 > 0$.

*Remark.* If $\beta_1(a) = [\beta_0]_1 a_1 e^{-\alpha_1 a_1}, \mu_1(a_1) \equiv \mu_1$ then the solution of (12) is

$$r_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_1]_1(P_1) = 0, [f_2](Q_1) = \lambda_1 Q_1\ h_1(Q_1) = 1, \qquad and\ \ g_1(P_1) = c_1 P_1$$

$$[\rho_1]_1 + [\rho_a]_1 = -(\mu_1 + \lambda_1 Q_1)\rho_1$$

$$\rho_1(t, 0) = \int_0^{\infty} [\beta_0]_1 a_1 e^{-\alpha_1 a_1} \rho_1(t, a_1)\, da_1$$

$$\frac{dQ_1}{dt} = -b_1 Q_1 + c_1 P_1 Q_1$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_1}{dt} = \omega_1 P_1 - \lambda_1 P_1 Q_1$$

$$\frac{dQ_1}{dt} = b_1 Q_1 - c_1 P_1 Q_1$$

Where $\omega_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$ THE EQUILIBRIUM $(\rho_0)_1(a_1), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_1(a_1), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_1 + (\rho_a)_1 = -\big(\mu_1(a_1) + (f_1)_1(P_1)\big)\rho,$$

$$\rho_1(t,0) = \int_0^\infty \beta_1(a_1)\rho_1(t,a_1)da_1 \qquad (13)$$

The solution of (13) has the form of

$$\rho_1(t,a_1) = Ce_1^{r(t-a_1)} - \int_0^a (\tau)d\tau - \int_0^1 (f_1)_1\big(P_1(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_1(t) = C_1 e^{r_1 t - \int_0^t (f_1)_1(P_1(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_1^{-1}(r)$ if $k_1 > r$.

This is equivalent to

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{(k_1)_1 a_1} da_1 < 1.$$

In fact, if we define

$$R_1(s_1) = \int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau)d\tau} e^{-s_1 a_1}\, da_1$$

$R(s)$ is a decreasing function. $R_1(r_1) = 1$ And $R_1(k_1)_1 < 1$ imply $r_1 < (k_1)_1$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_1(a_1), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{k_1 a_1} da_1 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_1 \equiv \int_0^\infty (\rho_0)_{11}(a_1)da_1$ *satisfies*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{-(f_1)_1(P_0)a_1} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_1(a_1), 0)$

Let $u_1(t, a_1) = \rho(t, a_1) - \rho_0(a_1)$ Then

$$N_1(t) \equiv \int_0^\infty u_1(t, a_1)da_1 = \int_0^\infty \rho_1(t, a_1)da_1 - \int_0^\infty (\rho_0)_1(a_1)da = P_1(t) - P_0)_1$$

And

$$(u_1)_1 + (u_a)_1 = (\rho_t)_1 + (\rho_a)_1 - (\rho_0^1)_1(a_1)$$

$$= -\mu_1(a_1) + (f_1)_1(P_1) + (f_2)_1(Q_1))(u_1 + (\rho_0)_1) + \mu_1(a_1)(\rho_0)_1 + (f_1)_1(P_0)_1(\rho_0)_1$$

$$= -\mu_1(a_1)u_1 - (f_1)_1(P_1) + (f_2)_1(Q_1))u_1 - (\rho_0)_1(f_1)_1(P_1) - (f_1)_1(P_0)_1$$

$$+ (f_2)_1(Q_1)), \qquad (15)$$

Where $(\rho_0^1)_1(a_1)$ denotes $d(\rho_0)_1(a_1)/da_1$

Expansion $(f_1)_1(P_1)$ and $(f_2)_1(Q_1)$ at $(P_0)_1$ and 0 yield respectively

$$(f_1)_1(P_1) = (f_1)_1((P_0)_1 + N_1)$$

$$= (f_1)_1(P_0)_1 + (f_1^1)_1(P_0)_1N_1 + o_1(N_1), \qquad (16)$$

$$(f_2)_1(Q_1)$$

$$= (f_2^1)_1(0)Q + o_1(Q_1) \qquad (17)$$

Write $f_1(P_0) = f_{10}, f_1'(P_0) = f_{10}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_1)_1 + (u_a)_1 \approx -(\mu_1(a_1) + (f_{10})_1)u - (f_{10}^1)_1(\rho_0)_1(N_1) - (f_{20}^1)_1(\rho_0)_1(Q_1)$$

$$u_1(t,0) = \rho_1(t,0)(\rho_0)_1(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t,a) + (\rho_0)_1(a_1)da_1 - (\rho_0)_1(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t,a_1)da_1 + (h_1(Q_1) - 1)(\rho_0)_1(0) \qquad (18)$$

Expanding

$$h_1(Q_1) = 1 + (h^1)_1(0)Q_1 + 0(Q_1) \equiv 1 + h_0'Q_1 + 0Q_1$$

And substituting it into (18), we arrive at

$$u_1(t,0) \approx \int_0^\infty \beta_1(a_1)u_1(t,a_1)da_1 + (h_{a_1}')(\rho_0)_1Q_1$$

The linearization is the following:

$$[u_1] + [u_1] = -(\mu_1(a_1) + [t_{10}]_1u_1 - [\alpha_1]_1(a_1)N_1 - [\alpha_2]_1(a_1)Q_1$$

$$u_1(t,0) = \int_0^\infty \beta_1(a_1)u_1(t,a_1)da_1 + [\xi_1]Q_1 \qquad (19)$$

$$\frac{dQ_1}{dt} = -b_1Q_1$$

where

$$[\alpha_1]_1(a_1) = [t_1^1 a]_1[p, o]_1(a_1)d2(a_1) = [f_{20}^1]_1(\rho\, o\,)_1(a_1)\xi = [\mathrm{h\ o}^1]_1[\rho \mathrm{o}]_1(\mathrm{o}) \text{ and } \mathrm{b}_1 = \mathrm{b}_1 - g_1(P_1)$$

If $g_1(P_0) > b$, clearly, $(\rho_0)_1(a_1), 0)$ is unstable and the predator population is persistent. Assume $g_1(P_0)_1 < b$.

Then, $Q_1(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_1 + u_a = -(\mu_1(a_1) + (f_{10})_1)u_1 - (\alpha_1)_1(a_1)N_1$$

$$= -(\mu_1(a_1) + f_{10})u_1 - \alpha_1(a_1)\int_0^\infty u_1(t, a_1)da_1$$

$$u_1(t, 0) = \int_0^\infty \beta_1\,(a_1)u_1(t, a_1)da_1$$

Using the method developed by Cushing [2, 3], we suppose

$$u_1(t, a_1) = \tilde{\mathrm{u}}_1(a_1)^{e^{\gamma(t-a_1)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_1}\, da_1$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_1)$ gives,

$$\tilde{\mathrm{u}}_1(a_1) = e^{-\int_0^a \mu_1(\tau_1)d\tau - (f_{10})_1 a}\left(\tilde{\mathrm{u}}_1(0) - w\int_0^a e^{\int_0^\tau \mu_1(s_1)ds + (f_{10})}\tau_1\,(\alpha_1)_1(\tau_1)e^{\gamma_1\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_1(a_1)e^{-\gamma_1 a_1 - (f_{10})_1 a_1 - \int_0^a \mu_1(\tau_1)d\tau}\right)\tilde{\mathrm{u}}_1(0)$$

$$+w\int_0^{\infty}\beta_1(a_1)e^{-\gamma_1 a_1-(f_{10})_1 a_1-\int_0^a \mu_1(\tau_1)d\tau}\int_0^a(\alpha_1)_1(\tau_1)e^{-\gamma_1\tau_1}e^{(f_{10})_1\tau_1+\int_0^{\tau}\mu_1(s_1)ds}d\tau_1\,da_1$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty}e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}da\right)u(0)$$

$$+\left(-1-\int_0^{\infty}e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}\int_0^a e^{f_{10}\tau+\int_0^{\tau}\mu(s)ds+y\tau}\alpha_1(\tau)d\tau\,da\right)w=0.\ (23)$$

Define $Y(\gamma,a)=e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta\,(a)Y(\gamma,a)da\right)u(0)$$

$$+w\int_0^{\infty}\beta(a)Y\,(\gamma,a)\int_0^a\alpha_1(\tau)Y^1(\gamma,\tau)d\tau\,da=0,(24)$$

$$\left(\int_0^{\infty}Y(\gamma,a)da\right)u(0)$$

$$+\left(-1-\int_0^{\infty}Y(\gamma,a)\int_0^a\alpha_1(\tau)Y^{-1}(\gamma,\tau)d\tau\,da\right)w=0.$$

Equation (20) has a solution of the form of $u(t,a)=u(a)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $u(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I(y,a)=\int_0^a\rho_0(\tau)Y^{-1}\,(\gamma,\tau)d\tau$. We obtain the characteristic equation

$$f_{10}^{l}\left(\int_0^{\infty} Y(\gamma,a)da\right)\left(\int_0^{\infty}\beta(a)Y(\gamma,a)I(\gamma,a)da\right)$$

$$+\left(1-\int_0^{\infty}\beta(a)Y(\gamma,a)da\right)\left(1+f_{10}^{l}\int_0^{\infty}Y(\gamma,a)I(\gamma,a)da\right)=0. \quad (25)$$

Since

$$\rho_0(a)=\rho_0(0)e^{-\int_0^a \mu(\tau)d\tau - f_{10}a}$$

$$I(\gamma,a)=\int_0^a \rho_0(\tau)e^{\gamma\tau+f_{10}\tau+\int_0^{\tau}\mu(s)ds}d\tau$$

$$=\frac{\rho_0(0)}{\gamma}(e^{\gamma a}-1)$$

Let

$$E(a)=e^{-f_{10}a-\int_0^a \mu(\tau)d\tau}$$

Then $[Y_1]_1[\gamma_1]_1,a_1)=\ E_1(a_1)e^{-\gamma_1 a_1}$. by a simple calculation,

$$\int_0^{\infty}\beta_1(a_1)Y_1(\gamma_1,a_1)I(\gamma_1,a_1)da=\frac{[\rho_0]_1(0)}{\gamma}\left(1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}\ da\right);$$

$$\int_0^{\infty}Y_1(\gamma_1,a_1)I(\gamma_1,a_1)da=\frac{[\rho_0]_1(0)}{\gamma_1}\left(\int_0^{\infty}E_1(a_1)e^{-\gamma_1 a_1}\ da_1\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_1(0)}{\gamma_1}[f_{10}^{1}]\int_0^{\infty}E_1(a_1)e^{-\gamma_1 a_1}da_1\ (\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1)$$

$$+\left(1-\int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right).$$

$$\left(1+\frac{(\rho_0)_1(0)}{\gamma_1}(f_{10}^{'})_1\left(\int_0^\infty E_1(a_1)da_1-\int_0^\infty E_1(a_1)e^{-\gamma_1 a_1}da_1\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right)(1+\frac{(f_{10}^{'})_1}{\gamma_1}(\rho_0)_1(0)\int_0^\infty E_1(a_1)da_1)=0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_1(a_1)Y_1(\gamma_1 a_1)da_1 = 1-\int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1 \neq 0,$$

Since $\int_0^\infty \beta_1(a_1)Y_1(\gamma_1 a_1)\int_0^a \alpha_1(\tau_1)Y^{-1}(\gamma_1,\tau_1)d\tau_1 da_1 > 0$. Therefore, the characteristic solution is

$$\gamma_1 = -(f_{10}^{'})_1(\rho_0)_1\int_0^\infty E_1(a_1)da_1.$$

Since $f_1'(P_0)_1 > 0$, we obtain the following result

THEOREM 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_1), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_1), 0)$ which is related to the extinction of debit summation as has been referred to in the foregoing and to The possible persistence of the prey

population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_1 + (\mu_1(a_1) + [f_1]_1(P_1^*) + [f_2]_1(Q_1^*))P_1^* = 0, \quad (26)$$

$$p_1^*(0) = h_1(Q_1^*)\int_0^\infty \beta_1(a_1)p_1^*(a_J)p_1^*(a_1)da_1 \quad (27)$$

$$g_1(P_1^*) - b_1 = 0, \quad (28)$$

Where $P_1^* = \int_0^\infty p_1^*(a_1)\, da_1$

Equation (28) yields $P^* = g^{-1}(b)$,

From (26),

$$p_1^*(a) = p_1^*(0)e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}$$

Substituting into (27), we need to solve

$$\left(1 - h_1(Q_1^*)\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}da_1\right)p_1^*(0) = 0$$

For $Q_1^*$

Fix $p_1^*(0)$ and $P_1^*$ Let

$$H_1(z_1) = 1 - h_1(z_1)\int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1}da_{,1}$$

And assume that

$$[k_3]_1 \int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(g_1^{-1}(b_1))a_1 - [k_2]_1 a_1}\, da_1 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_1(z_1)] > 0.$$

Clearly,

$$\frac{dH_1}{dz_1} = h_1(z)(\,[f_2^{'}]_1(z_1)\,) \int_0^{\infty} a_1\beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - f_1(P_1^*)a_1 - f_2(z_1)a_1} da_1$$

$-h^{'}(z_1)\int_0^{\infty} \beta_1(a_1)\, e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1}\ da_1$

By the assumptions on $\left[f_2^1\right]_1 and\ [h^1]^1 dH_1/dz > 0$. Hence, if

$$H_1(0) = 1 - \int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1} da_1 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_1(0) < 0$ follows if $g_1^{-}(b_1) < [f_1^{-1}]_1(r_1)$, where $r_1$ is the solution of (12).

Indeed, since

$$R_1(s_1) = \int_0^{\infty} \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 e^{-y_1^a}} da_1$$

Is a decreasing function of $s_1$ when $P_1^* = [g^{-1}]_1(b_1) < [f_1^{-1}]_1(r_1)$, that is $[f_1]_1(P_1^*) < r$, $R\big(f_1(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_1(g^{-1}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1\left(g_1^{-1}(b)\right)a_1} > 1. \qquad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_1^*(a_1), Q_1^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_1^*(a_1), Q_1^*)$ is more complicated. We again consider perturbations $m_1(t_1 a_1) = p_1(t_1, a_1) - p_1^*(a_1)$ and $W_1(t) = Q_1(t) - Q^*$.Then

$$M_1(t) \equiv \int_0^\infty m_1(t, a_1)da_1 = P_1(t) - P_1^*,$$

$$[m_1]_1 + [m_a]_1 = [p_1]_1 + [p_a]_1 - \left[p^*\right]_1(a_1) \approx -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1$$

$$-[f_{1*}^1]_1 p_1^*(a_1)M - [f_{2*}^1]_1 p_1^*(\mathrm{a})W_1,$$

$$m_1(t,0)p_1(t,0) - p_1^*(0) \approx h_* \int_0^\infty \beta_1(a_1)m_1(t,a_1)\,da + \frac{[h_*^1]_1}{[h_*]_1}\, p_1^*(0)W_1$$

$$\frac{dW_1}{dt} = \frac{dQ_1}{dt} \approx [g_*^1]_1 Q_1^* M_{,1}$$

Where $[f_{1*}]_1$ denoting $[f_1(P^*)]_1$, $[f_{1*}]_1$ denoting $[f_2(P^*)]_1$, $\left[f_{1*}^{'}\right]_1$ denoting $\left[f^{'}{}_1(P^*)\right]_1 \left[f_{2*}^{'}\right]_1$ denoting

$[f_2]_1\,(Q_1^*,\quad h_{1*}^{'}$ denoting $h_*^{'}$ denoting $h_1^{'}\,(Q_1^*)$, and $\left[g_*^{'}\right]_1$

Denoting $\left[g^{'}\right]_1(P_1^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_1]_1 + [m_a]_1 = -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1 - [f'_{1*}]_1[p_1^*](a_1)M_1 - [f'_{2*}]_1 p_1^*(a_1)W_1$$

$$m_1(t,0) = [h_*]_1 \int_0^\infty \beta_1(a_1)m_1(t,a_1)da_1 + \frac{[h_*^1]_1}{[h^*]_1}[p^*]_1(0)\, W_{,1} \qquad (30)$$

$$\frac{dW_1}{dt} = [g_*^1]_1[Q_1^*]M_1$$

Integrating $[m_1]_1(t,a_1)$ with respect to a from 0 to ∞, we have

$$\frac{dM_1(t)}{dt} = \int_0^\infty [mt]_1(t,a_1)da_1$$

$$= -\int_0^\infty [m_a]_1\,(t,a_1)da_1 - \int_0^\infty \mu(a_1)m(t,a_1)da_1$$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f'_{1*}]_1 \int_0^\infty p_1^*(a_1)da_1\right)M_1 - [f_{2*}^1]_1 \int_0^\infty p_1^*(a_1)da_1 W_1$$

$$= m_1(t,0) - \int_0^\infty \mu_1(a_1)\, m_1(t,a_1)da_1$$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f_{1*}']_1 P_1^*\right)M_1 - [f_{2*}']_1 P_1^* W_{,1}$$

From (30), we arrive at the equations

$$\frac{dW_1(t)}{dt} = \int_0^\infty ([h_*]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1$$

$$-\left([f_{1*}] + [f_{2*}]_1 + [f_{1*}']_1 P_1^*\right)M_1 + \left(\frac{[h_*']_1}{[h_*]^1}\, P_1^*(0) - f_{2*}' P_1^*\right)W_1,$$

(31)

$$\frac{dW_1(t)}{dt}\left[g^{'}{}_{*}\right]_1 Q_1^{*} M_{.1}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{*}]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1 = ([h_{*}]_1\beta_1(\theta_1) - \mu_1(\theta_1))M_{,1}$$

Where $0< \theta_1 = \theta_1(t) < \infty$.

And if $a_1^2\,[a_1^2]_1<4\,[a_2]_1<4\,a_2 e$,

Re $\lambda_1\left[\lambda_1\right]_1 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_1 = -\frac{1}{2}a_1[a2\ ]_1 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_{*}\,[h_{*}]_1\beta_1(a_1) - \mu_1(a_1)\right|$. Then$|n_1(t)|$ ⊠Δ∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium($\rho^{*}(a), \boldsymbol{Q}^{*}$) is asymptotically stable if either $a_1^2 <4\,a_2 e$ and $[a_1] > 2\Delta$,or $[a_1^2]$ ⊠4$[a_2]$e and $(a_1 - \sqrt{(a_1^2)_1 - 4(a_2)_1 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+1}\right)$ and $\beta_1(a_1) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1) + d\tau_1 da} = \frac{1}{(m+\mu_0+1)^2} + \frac{2}{(m+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_1 + [\mu_0]$ 0.5. Assume$[f_1](x_1) = K_1 x_1$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_1(x_1) = [0x_1]/(1+x_1)$; the functional responses $[[f_2](x) = [0]\, x_1/(1+x_1)$and $h_1(x_1)$=(1+$\omega_1 x_1$)/(1 + $x_1$)(See Holling [7] and Rashevsky [13])

Equation (14),now has the form of

$$\frac{m_1 + [\mu_0] + [K][P_0])}{[m_1]_1 + (\mu_0)_1 + [K][P_0])_1^3\ [K]_1 [P_0]_1^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1+P_0) < b_1$i e $[P_0] < b/(\delta_1 - b_1)$then the equilibrium$(P_0)$ $[P_0]_1(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_1$

$[f_{1*}] = K_1\ P^*\ [f_{2*}] = \frac{\sigma Q_1^*}{1+Q_1^*} \quad [h_*] = \frac{1+\omega_1 Q_1^*}{1+Q_1^*}$

$[f_{1*}] = K_1\ [f_{2*}^1] = \frac{\sigma_1}{(1+Q_1^*)^2} [\ g_*^1] = \frac{\sigma_1}{(1+P_1^*)^2} \quad [h_*^1] = \frac{\omega_1 - 1}{(1+Q_1^*)^2}$

## AN ANALOGY BETWEEN PREDATOR PREY AND ACCENTUATION ATTRION MODELS

**Best example how the model could be applied is the Schrödinger's equation itself wherein RHS gobbles up LHS, and the terms on RHS do the same with each other. It is like adding milk to water. Whether it happens instantaneously or time lag is there is a different question altogether dealt at the opportune points**

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| **Best example that can be given is the formulation of the quantum wave equation albeit in a nascent form for the recently found gravitational waves. This has been studied in length in various papers of gravitational waves, wherein the attribution of how the model is used is described. Note the calculation holds for only one module. This is in consideration to the fact that the relation is beteen the both debit systems and credit systems which as we have enucleated in the foregoing bears ample testimony for the application of the model** |
| **Gravitational quantum wave equation:** |
| **Equation Time-dependent equation (Wikipedia)** |
| **(1) In quantum mechanics, the Schrödinger equation is (=) a partial differential equation that describes how the quantum state of a quantum system changes with (e&eb) time. It was formulated in late 1925, and published in 1926, by the Austrian physicist Erwin Schrödinger.[1]** |
| **(2) In classical mechanics Newton's second law, (F = ma), is used to make (eb) a mathematical prediction as to what path a given system will take following a set of known initial conditions.** |
| **(3) In quantum mechanics, the analogue of Newton's law is (=) Schrödinger's equation for (e) a quantum system (usually atoms, molecules, and subatomic particles whether free, bound, or localised).** |
| **(4) It is not a simple algebraic equation, but in general a linear partial differential equation, describing (eb) the time-evolution of the system's wave function (also called a "state function").[2]:1–2** |
| **(5) The concept of a wavefunction is (=) a fundamental postulate of quantum mechanics.** |
| **(6) Although Schrödinger's equation is often presented as a separate postulate, some authors [3]: Chapter 3 show that some properties resulting from Schrödinger's equation may be deduced just from (e) symmetry principles alone, for example the commutation relations.** |
| **(7) Generally, "derivations" of the Schrödinger equation demonstrate (eb) its mathematical plausibility for describing (eb) wave–particle duality, but to date there are no universally accepted derivations of (e) Schrödinger's equation from (e) appropriate axioms.** |
| **(8) In the Copenhagen interpretation of quantum mechanics, the wave function is (=) the most complete description that can be given of (e) a physical system.** |
| **(9) Solutions to Schrödinger's equation describe (eb) not only molecular, atomic, and subatomic systems, but also macroscopic systems, possibly even the whole universe.[4]:292ff** |

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| **(10) The Schrödinger equation is consistent with (e&eb) classical mechanics but not with (e&eb) special relativity.** |
| **(11) Making quantum mechanics consistent with (e&eb) special relativity is one of the great achievements of (e) Quantum Field Theory and requires (e) the introduction of the creation and annihilation operators so that (e) the number of particles can no longer be considered constant.** |
| **(12) The Schrödinger equation is not the only way to make predictions in (eb) quantum mechanics—other formulations can be used, such as Werner Heisenberg's matrix mechanics, and Richard Feynman's path integral formulation** |
| **(13) The form of the Schrödinger equation depends on (=) the physical situation (see below for special cases).** |
| **(14) The most general form is the time-dependent Schrödinger equation, which gives a description of a system evolving with time: [5]:143** |
| **(15) A wave function that satisfies the non-relativistic Schrödinger equation with V = 0. In other words, this corresponds to (e &eb) a particle traveling freely through (e&eb) empty space. Time-dependent** |
| **Schrödinger equation (general)** |
| **(16) i ℏ ∂ ∂ t Ψ ( r , t ) = H ^ Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t) where i is the imaginary unit, ħ is the Planck constant divided by 2π, the symbol ∂/∂t indicates a partial derivative with respect to time t, Ψ (the Greek letter psi) is the wave function of (e) the quantum system, and Ĥ is the Hamiltonian operator (which characterizes the total energy of any given wave function and takes different forms depending on the situation).** |
| **(17) Each of these three rows is a wave function which satisfies (eb) the time-dependent Schrödinger equation for a harmonic oscillator. Left: The real part (blue) and imaginary part (red) of the wave function. Right:** |
| **(18) The probability distribution of finding the particle with (e&eb) this wave function at a given position.** |
| **(19) The top two rows are examples of stationary states, which correspond to (e&eb) standing waves.** |
| **(20) The bottom row is an example of a state which is (=) not a stationary state.** |
| **(21) The right column illustrates (eb) why stationary states are (=) called "stationary".** |

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| **(22) The most famous example is the non-relativistic Schrödinger equation for (e) a single particle moving in an electric field (but not a magnetic field; see the Pauli equation): [6]** |
| **Time-dependent Schrödinger equation** |
| **(Single non-relativistic particle)** |
| **(23) i ℏ ∂ ∂ t Ψ ( r , t ) = [ − ℏ 2 2 μ ∇ 2 + V ( r , t ) ] Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)** |
| **(24) Here μ is (=) the particle's "reduced mass", V is its potential energy, ∇2 is the Laplacian (a differential operator), and Ψ is (=) the wave function (more precisely, in this context, it is called the (=) "position-space wave function").** |
| **(25) In plain language, it means "total energy equals (=) kinetic energy plus potential energy", but the terms take unfamiliar forms for (e) reasons explained below.** |
| **(26) Given the particular differential operators involved, this is (=) a linear partial differential equation.** |
| **(27) It is also a diffusion equation, but unlike the heat equation, this one is (=) also a wave equation given (e&eb) the imaginary unit present in the transient term.** |
| **(28) The term "Schrödinger equation" can refer to (e) both the general equation (first box above), or the specific nonrelativistic version (second box above and variations thereof).** |
| **(29) The general equation is indeed quite general, used throughout quantum mechanics, for everything from the Dirac equation to quantum field theory, by plugging in various complicated expressions for the Hamiltonian. The specific nonrelativistic version is a simplified approximation to reality, which is quite accurate in many situations, but very inaccurate in others (see relativistic quantum mechanics and relativistic quantum field theory)** |
| **(30) To apply the Schrödinger equation, the Hamiltonian operator is set up for the system, accounting for the kinetic and potential energy of the particles constituting the system, then inserted into the Schrödinger equation. The resulting partial differential equation is solved for the wave function, which contains information about the system.** |
| **Time-independent equation** |

**(31) The time-dependent Schrödinger equation described above predicts that wave functions can form standing waves, called stationary states (also called "orbitals", as in atomic orbitals or molecular orbitals). These states are important in their own right, and if the stationary states are classified and understood, then it becomes easier to solve the time-dependent Schrödinger equation for any state. Stationary states can also be described by a simpler form of the Schrödinger equation, the time-independent Schrödinger equation. (This is only used when the Hamiltonian itself is not dependent on time explicitly. However, even in this case the total wave function still has a time dependency.)**

**Time-independent Schrödinger equation (general)**

**(32) E Ψ = H ^ ⊠ Ψ {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi } {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi }**

**In words, the equation states:**

**(33) When the Hamiltonian operator acts on a certain wave function Ψ, and the result is proportional to the same wave function Ψ, then Ψ is a stationary state, and the proportionality constant, E, is the energy of the state Ψ.**

**(34) The time-independent Schrödinger equation is discussed further below. In linear algebra terminology, this equation is an eigenvalue equation.**

**(35) As before, the most famous manifestation is the non-relativistic Schrödinger equation for a single particle moving in an electric field (but not a magnetic field):**

**Time-independent Schrödinger equation (single non-relativistic particle)**

**(36) E Ψ ( r ) = [ − ℏ 2 2 μ ∇ 2 + V ( r ) ] Ψ ( r ) {\displaystyle E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )} E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )**

**Gravitational waves: (Wikipedia)**

**(37) In theory, the loss of energy through gravitational radiation could eventually drop the Earth into the Sun. However, the total energy of the Earth orbiting the Sun (kinetic energy + gravitational potential energy) is about 1.14×1036 joules of which only 200 Watts (joules per second) is lost through gravitational radiation, leading to a decay in the orbit by about 1×10−15 meters per day or roughly the diameter of a proton. At this rate, it would take the Earth approximately 1×1013 times more than the current age of the Universe to spiral onto the Sun. This estimate overlooks the decrease in r over time, but the majority of the time the bodies are far apart and only radiating slowly, so the difference is unimportant in this example.**

**More generally, the rate of orbital decay can be approximated by[39]**

**(38)** $\frac{dr}{dt} = -\frac{64}{5}\frac{G^3}{c^5}\frac{(m_1 m_2)(m_1 + m_2)}{r^3}$ , {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,} {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,}

**where r is the separation between the bodies, t time, G Newton's constant, c the speed of light, and m1 and m2 the masses of the bodies. This leads to an expected time to merger of [39]**

$t = \frac{5}{256}\frac{c^5}{G^3}\frac{r^4}{(m_1 m_2)(m_1 + m_2)}$. {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.} {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.}

**(39) For example, a pair of solar mass neutron stars in a circular orbit at a separation of 1.89×108 m (189,000 km) has an orbital period of 1,000 seconds, and an expected lifetime of 1.30×1013 seconds or about 414,000 years. Such a system could be observed by LISA if it were not too far away. A far greater number of white dwarf binaries exist with orbital periods in this range. White dwarf binaries have masses in the order of the Sun, and diameters in the order of the Earth. They cannot get much closer together than 10,000 km before they will merge and explode in a supernova which would also end the emission of gravitational waves. Until then, their gravitational radiation would be comparable to that of a neutron star binary.**

**(40) By substituting value of dr/dt in Schrödinger's wave equation we get the quantum wave equation for gravitational waves.**

In order to have (H3) and (H4) hold, $m, \mu_0$, K, b, δ, and $\omega$ need to satisfy

$$\frac{\omega_1}{(m + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 + b_1} + \sigma)^2} + \frac{2\omega_1}{([m_1]_1 + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 + b_1} + \sigma_1)^3} < 1$$

And

$$\frac{1}{(m_1 + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 - b_1})^2} + \frac{2\omega_1}{(m_1 + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 - b_1})^3} > 1$$

When the inequalities (33) and (34) hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$(m_1 + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 - b_1} + (m_1 + [\mu_0]_1 + \frac{K_1 b_1}{\delta_1 - b_1} + \sigma_1)\, Q_1^*)^3$

$= m_1 + [\mu_0]_1 + 3 + \frac{K_1 b_1}{\delta_1 - b_1}(1 + Q^*)(1 + \omega_1 Q_1^*$

Now

$$\left[a_1^2\right] - 4\left[a_2\right] e = \left(\frac{2K_1 b_1}{\delta_1 - b_1}\right)^2 + \frac{4K_1 b_1 \sigma_1 Q_1^*}{(\delta_1 - b_1)(1 + Q_1^*)} + \frac{\sigma_1^2 Q_1^{*2}}{(1 + Q_1^*)^2}$$

$$\frac{4\, b_1 \sigma_1 (\delta_1 - b) Q^*}{\delta_1 (1 + Q_1^*)^2} - \frac{4\, (\delta_1 - b_1)^2 (1 - \omega_1) Q_1^*}{\delta_1 (1 + \omega_1 Q_1^*)(1 + Q_1^*)} P^*(0)$$

And

$\Delta = \lim_{0 < a < \infty} |[h_*]_1\, a e_1^{mu}\ [\mu_0]_1 - \frac{a_1}{a_1 + 1}$

Hence if $[a_1^2 > 4(a_2)_1 e$ and

$$[a_1] = \frac{2K_1 b_1}{\delta_1 - b_1} + \frac{\sigma_1 Q_1^*}{1 + Q_1^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_1]_1 = [f_{1*}]_1 + [f_{2*}]_1 + [f_{1*}] P_1^*$$

$$[a_2]_1 = \left[f^{1}_{2*}\right] P^{*}_{1} - \frac{h^{1}_{*}}{h_{*}} P^{*}(0)$$

$$e = [g^{1}_{*}]_1 Q^{*}_{1}$$

$$\eta_1(t) = \left[h_{*}\right]_1 \beta_1 \left(\theta_1(t)\right) - \mu_1(\theta(t)),$$

The system(31)can be written as the following system:

$$[M(t)]\left[a_1\right]M_1 - [a_2]_1 W_1 - [n(t)]M_1$$

$$w_1(t) = eM_1$$

Write

$$A = \begin{pmatrix} -[a_1]_1 - [a_2]_1 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)= \begin{pmatrix} [\eta(t)]_1 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t) = \begin{pmatrix} M_1 \\ W_1 \end{pmatrix}$$

The solution of (32)is

$$X_1(t) = e^{A,t}[X]_1(0) + \int_0 e^{At}$$

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# Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

## (Module Three)

**Abstract: Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-435) 1990. We demonstrate its applicability cosmology/ quantum computation economics et al. As stated in the foregoing we shall apply model to various cosmological problem and show the relevance to the same. . To exemplify model can be applied to the following situational significations: Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an**

**outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 Masaru Shibata and Takashi Nakamura** perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. **Authors** use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society **Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995**

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(g_3)Q_3$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_3(0) = 0, \quad g_3(P_3) > 0, \quad P_3 \text{ 0}, \qquad \lim_{P_3 \to \infty} g_3(P_3) = (K_3)_3 \text{ ∞}. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_3 + (\rho_a)_3 + D_3(a_3, P_3, Q_3)\rho_3 + 0, \qquad 0 < a_3 \text{ ∞},$$

$$\rho_3(t,0) = \int_0^\infty n_3(a_3, P_3, Q_3)\rho_3(t_3, a_3)da_3$$

Where $D_3$ is the mortality function and $n_3$ is the fertility function.

Put $_3$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_3}{dt}\right]_3 = x_3\left[\frac{dQ_2}{dt}\right]_3 - x_2\left[\frac{dQ_3}{dt}\right]_3$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_3 - x_4\left[\frac{dQ_2}{dt}\right]_3$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_3 - x_6\left[\frac{dQ_3}{dt}\right]_3$$

Where $x_3, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_3$ and $g_3(P_3)Q_3$

Now when you add the term $(b_3Q_3)$ with $g_3(P_3)(Q_3)$ we are infect see is that $(b_3Q_3)$ is credited incrementally part of $g_3(P_3)(Q_3)$ or instant nearly, when subtracted from $\frac{dQ_3}{dt}$ , it means that $g_3(P_3)(Q_3)$ is slowly eaten by $\frac{dQ_3}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_3 = \mu_3(a_3) + (f_3)_3(P_3)_3 + (f_2)_3(Q)_3 \quad (2)$$

$$n_3 = \beta_3(a_3)h_3(Q_3), \quad (3)$$

Where $\mu_3$the age- dependent debit rate whit is $\mu_3(a_3)$ 0 $\forall a > 0$; $(f_3)_3$ is the density- dependent death rate with $f_3(0) = 0$, $(f_3^3)_3(P_3)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_3)_3(P_3) = (k_3)_3$ ; $f_2$ is the functional response of

creditors to prey with $[f_2]_3(0) = 0, (f_2^3)_3(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_3\ (Q_3) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 3, h^3(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_3 + (\rho_a)_3 = -\mu_3(a_3) + (f_3)_3 P_3) + \big(f_2)_3(Q_3)\big)(\rho_3)_3 \tag{4}$$

$$\rho_3(t_3, 0) = h_3(Q_3)\int_0^\infty \beta_3\,(a_3)\rho_3(t, a_3)da, \tag{5}$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(P_3)Q_3 \tag{6}$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_3$ and $Q_3$ of (4-6) satisfy

$$0 \ \rho_3(t, a_3) \ \rho_3^*(t, a_3),$$

$$0 \ Q_3(t) \ Q_3^*(t),$$

Where $\rho_3^*(t, a_3)$ and $Q_3^*(t)$, are solutions of

$$(\rho_t^*)_3 + (\rho_a^*)_3 = -(\mu_3(a_3) + (f_3)_3(P_3^*) + (f_2)_3 Q_3^*))\,\rho_3^*, \tag{7}$$

$$\rho_3^*\,(t, 0) = \int_0^\infty \beta_3(a_3)\,\rho_3^*\,(t, a_3)da_3 \tag{8}$$

$$\frac{dQ_3^*}{dt} = -\,b_3\,Q_3^* + g_3(P_3^*)Q_3^*, \tag{9}$$

With $P_3^*(t) = \int_0^\infty \rho_3^*\ (t, a_3)da_3$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_3^*(t)}{dt} = P_3^*(t)\big(r_3 - (f_3)(P_3^*) - (f_2)_3(Q^*)\big), \quad (10)$$

$$\frac{dQ_3^*(t)}{dt} = Q_3^*(t)(-b_3 + g_3(P_3^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_3(a)e^{-\int_0^a \mu_3(\tau)dt}e^{-ra}\ da_3 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_3^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(t)dt}\ da_3 > 1, \quad (H1)$$

Which guarantees $r_3 > 0$.

*Remark.* If $\beta_3(a) = [\beta_0]_3 a_3 e^{-\alpha_3 a_3}, \mu_3(a_3) \equiv \mu_3$ then the solution of (12) is

$$r_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_3]_3(P_3) = 0, [f_2](Q_3) = \lambda_3 Q_3\ h_3(Q_3) = 1, \quad and\ \ g_3(P_3) = c_3 P_3$$

$$[\rho_3]_3 + [\rho_a]_3 = -(\mu_3 + \lambda_3 Q_3)\rho_3$$

$$\rho_3(t,0) = \int_0^{\infty} [\beta_0]_3 a_3 e^{-\alpha_3 a_3} \rho_3(t,a_3)\, da_3$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + c_3 P_3 Q_3$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_3}{dt} = \omega_3 P_3 - \lambda_3 P_3 Q_3$$

$$\frac{dQ_3}{dt} = b_3 Q_3 - c_3 P_3 Q_3$$

Where $\omega_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$

## THE EQUILIBRIUM $(\rho_0)_3(a_3), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_3(a_3), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_3 + (\rho_a)_3 = -\big(\mu_3(a_3) + (f_3)_3(P_3)\big)\rho,$$

$$\rho_3(t,0)$$

$$= \int_0^\infty \beta_3(a_3)\rho_3(t,a_3)da_3 \qquad (13)$$

The solution of (13) has the form of

$$\rho_3(t,a_3) = Ce_3^{r(t-a_3)} - \int_0^a (\tau)d\tau - \int_0^3 (f_3)_3\big(P_3(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_3(t) = C_3 e^{r_3 t - \int_0^t (f_3)_3(P_3(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_3^{-3}(r)$ if $k_3 > r$. This is equivalent to

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{(k_3)_3 a_3} da_3 < 1.$$

In fact, if we define

$$R_3(s_3) = \int_0^\infty \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau)d\tau} e^{-s_3 a_3}\, da_3$$

$R(s)$ is a decreasing function. $R_3(r_3) = 1$ And $R_3(k_3)_3 < 1$ imply $r_3 < (k_3)_3$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_3(a_3), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{k_3 a_3} da_3$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_3 \equiv \int_0^\infty (\rho_0)_{33}(a_3)da_3$ *satisfies*

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{-(f_3)_3(P_0)a_3} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_3(a_3), 0)$

Let $u_3(t, a_3) = \rho(t, a_3) - \rho_0(a_3)$ Then

$$N_3(t) \equiv \int_0^{\infty} u_3(t, a_3) da_3 = \int_0^{\infty} \rho_3(t, a_3) da_3 - \int_0^{\infty} (\rho_0)_3(a_3) da = P_3(t) - P_0)_3$$

And

$$(u_3)_3 + (u_a)_3 = (\rho_t)_3 + (\rho_a)_3 - (\rho_0^3)_3(a_3)$$

$$= -\mu_3(a_3) + (f_3)_3(P_3) + (f_2)_3(Q_3))(u_3 + (\rho_0)_3) + \mu_3(a_3)(\rho_0)_3 + (f_3)_3(P_0)_3(\rho_0)_3$$

$$= -\mu_3(a_3)u_3 - (f_3)_3(P_3) + (f_2)_3(Q_3))\, u_3 - (\rho_0)_3\, (f_3)_3(P_3) - (f_3)_3(P_0)_3$$

$$+ (f_2)_3(Q_3)), \qquad (15)$$

Where $(\rho_0^3)_3(a_3)$ denotes $d(\rho_0)_3(a_3)/da_3$

Expansion $(f_3)_3(P_3)$ and $(f_2)_3(Q_3)$ at $(P_0)_3$ and 0 yield respectively

$$(f_3)_3(P_3) = (f_3)_3((P_0)_3 + N_3)$$

$$= (f_3)_3(P_0)_3 + (f_3^3)_3(P_0)_3 N_3 + o_3(N_3), \qquad (16)$$

$$(f_2)_3(Q_3)$$

$$= (f_2^3)_3(0)Q + o_3(Q_3) \qquad (17)$$

Write $f_3(P_0) = f_{30}, f_3'(P_0) = f_{30}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_3)_3 + (u_a)_3 \approx -(\mu_3(a_3) + (f_{30})_3)u - (f_{30}^3)_3(\rho_0)_3(N_3) - (f_{20}^3)_3(\rho_0)_3(Q_3)$$

$$u_3(t, 0) = \rho_3(t, 0)(\rho_0)_3(0)$$

$$= h_3(Q_3)\int_0^{\infty} \beta_3(a_3)(u_3(t,a) + (\rho_0)_3(a_3)da_3 - (\rho_0)_3\,(0)$$

$$= h_3(Q_3)\int_0^{\infty} \beta_3(a_3)(u_3(t,a_3)da_3 + (h_3(Q_3) - 3)(\rho_0)_3\,(0) \qquad (18)$$

Expanding

$$h_3(Q_3) = 1 + \ (h^3)_3(0)Q_3 + 0(Q_3) \equiv 1 + h_0^{'}Q_3 + 0Q_3$$

And substituting it into (38), we arrive at

$$u_3(t,0) \approx \int_0^{\infty} \beta_3\ (a_3)u_3(t,a_3)da_3 + (h_{a_3}^{'})(\rho_0)_3Q_3$$

The linearization is the following :

$$[u_3] + [u_3] = -(\mu_3(a_3) + [t_{30}]_3u_3 - [\alpha_3]_3(a_3)N_3 - [\alpha_2]_3(a_3)Q_3$$

$$u_3(t,0) = \int_0^{\infty} \beta_3(a_3)u_3(t,a_3)da_3 + [\xi_3]Q_3 \qquad (19)$$

$$\frac{dQ_3}{dt} = -b_3Q_3$$

where

$$[\alpha_3]_3(a_3) = [t_3^3a]_3[p,o]_3(a_3)d2(a_3) = [f_{20}^3]_3(\rho\ o\ )_3(a_3)\xi = [\mathrm{h\ o}^3]_3[\rho\mathrm{o}]_3(\mathrm{o}) \text{ and } \mathrm{b}_3 = \mathrm{b}_3 - g_3(P_3)$$

If $g_3(P_0) > b$, clearly, $(\rho_0)_3(a_3), 0)$ is unstable and the predator population is persistent. Assume $g_3(P_0)_3 < b$.

Then, $Q_3(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_3 + u_a = -(\mu_3(a_3) + (f_{30})_3)u_3 - (\alpha_3)_3(a_3)N_3$$

$$= -(\mu_3(a_3) + f_{30})u_3 - \alpha_3(a_3)\int_0^\infty u_3(t, a_3)da_3$$

$$u_3(t, 0) = \int_0^\infty \beta_3\,(a_3)u_3(t, a_3)da_3$$

Using the method developed by Cushing [2, 3], we suppose

$$u_3(t, a_3) = \tilde{u}_3(a_3)^{e^{\gamma(t-a_3)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_3}\,da_3$. Then,

A solution of the first differential equation for $\tilde{u}(a_3)$ gives,

$$\tilde{u}_3(a_3) = e^{-\int_0^a \mu_3(\tau_3)d\tau - (f_{30})_3 a}\left(\tilde{u}_3(0) - w\int_0^a e^{\int_0^\tau \mu_3(s_3)ds + (f_{30})}\tau_3\,(\alpha_3)_3(\tau_3)e^{\gamma_3\tau}d\tau\right).$$

Substituting the expression (13) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\right)\tilde{u}_3(0)$$

$$+w\int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\int_0^a (\alpha_3)_3(\tau_3)e^{-\gamma_3\tau_3}e^{(f_{30})_3\tau_3 + \int_0^\tau \mu_3(s_3)ds}d\tau_3\,da_3$$

$$= 0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(\tau_3)d\tau} da\right)(\tilde{u}_3(0))_3$$

$$+\left(-3-\int_0^\infty e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(\tau_3)d\tau}\int_0^a e^{(f_{30})_3\tau_3+\int_0^\tau \mu_3(s_3)ds+y\tau}\alpha_3(\tau)d\tau\, da\right)w=0. \quad (23)$$

Define $Y(\gamma,a)=e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(t_3)d\tau}$. Then

$$\left(1-\int_0^\infty \beta_3\,(a_3)Y_3(\gamma_3,a_3)da_3\right)u_3(0)$$

$$+w\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)\int_0^a(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\, da_3=0, \quad (24)$$

$$\left(\int_0^\infty Y_3(\gamma_3,a_3)da_3\right)\tilde{u}_3(0)$$

$$+\left(-1-\int_0^\infty Y_3(\gamma_3,a_3)\int_0^a(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\, da_3\right)w=0.$$

Equqtion (20) has a solution of the form of $u_3(t,a_3)=u_3(a_3)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_3(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_3(\gamma_3,a_3)=\int_0^a(\rho_0)_3(\tau_3)Y^{-3}\,(\gamma_3,a_3)d\tau_3$. We obtain the characteristic equation

$$f_{30_3}^l\left(\int_0^\infty Y_3(\gamma_3,a_3)da_3\right)\left(\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$+\left(1-\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)da\right)\left(1+(f_{30}^l)_3\int_0^\infty Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_3(a_3)=(\rho_0)_3(0)e^{-\int_0^a \mu_3(\tau_3)d\tau_3-(f_{30})_3a_3}$$

$$I_3(\gamma_3,a_3)=\int_0^a (\rho_0)_3(\tau_3)e^{\gamma_3\tau_3+(f_{30})_3\tau+\int_0^\tau \mu_3(s_3)ds_3}d\tau$$

$$=\frac{(\rho_0)_3(0)}{\gamma_3}(e^{\gamma_3a_3}-1)$$

Let

$$E_3(a_3)=e^{-(f_{30})_3a_3-\int_0^a \mu_3(\tau_3)d\tau_3}$$

Then $[Y_3]_3[\gamma_3]_3,a_3)=E_3(a_3)e^{-\gamma_3a_3}$. by a simple calculation,

$$\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)I(\gamma_3,a_3)da=\frac{[\rho_0]_3(0)}{\gamma}\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3a_3}\ da\right);$$

$$\int_0^\infty Y_3(\gamma_3,a_3)I(\gamma_3,a_3)da=\frac{[\rho_0]_3(0)}{\gamma_3}\left(\int_0^\infty E_3(a_3)e^{-\gamma_3a_3}\ da_3\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_3(0)}{\gamma_3}[f_{30}^3]\int_0^\infty E_3(a_3)e^{-\gamma_3a_3}da_3\ (\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3a_3}da_3)$$

$$+\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right).$$

$$\left(1+\frac{(\rho_0)_3(0)}{\gamma_3}(f_{30}^{'})_3\left(\int_0^\infty E_3(a_3)da_3-\int_0^\infty E_3(a_3)e^{-\gamma_3 a_3}da_3\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right)(1+\frac{(f_{30}^{'})_3}{\gamma_3}(\rho_0)_3(0)\int_0^\infty E_3(a_3)da_3)=0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_3(a_3)Y_3(\gamma_3 a_3)da_3 = 1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3 \neq 0,$$

Since $\int_0^\infty \beta_3(a_3)Y_3(\gamma_3 a_3)\int_0^a \alpha_3(\tau_3)Y^{-3}(\gamma_3,\tau_3)d\tau_3 da_3 > 0$. Therefore, the characteristic solution is

$$\gamma_3 = -(f_{30}^{'})_3(\rho_0)_3\int_0^\infty E_3(a_3)da_3.$$

Since $f_3'(P_0)_3 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_3), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_3), 0)$ which is related to the extinction of debit summation types

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_3 + (\mu_3(a_3) + [f_3]_3(P_3^*) + [f_2]_3(Q_3^*))P_3^* = 0, \quad (26)$$

$$p_3^*(0) = h_3(Q_3^*)\int_0^\infty \beta_3(a_3)p_3^*(a_J)p_3^*(a_3)da_3 \quad (27)$$

$$g_3(P_3^*) - b_3 = 0, \quad (28)$$

Where $P_3^* = \int_0^\infty p_3^*(a_3)\, da_3$

Equation (28) yields $P^* = g^{-3}(b)$,

From (26),

$$p_3^*(a) = p_3^*(0)e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}$$

Substituting into (27), we need to solve

$$\left(1 - h_3(Q_3^*)\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}da_3\right)p_3^*(0) = 0$$

For $Q_3^*$

Fix $p_3^*(0)$ and $P_3^*$ Let

$$H_3(z_3) = 1 - h_3(z_3)\int_0^\infty \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3} da_{,3}$$

And assume that

$$[k_3]_3 \int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(g_3^{-3}(b_3))a_3 - [k_2]_3 a_3}\, da_3 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_3(z_3)] > 0.$$

Clearly,

$$\frac{dH_3}{dz_3} = h_3(z)([f_2^{'}]_3(z_3))\int_0^\infty a_3\beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - f_3(P_3^*)a_3 - f_2(z_3)a_3} da_{3}$$

$-h^{'}(z_3)\int_0^\infty \beta_3(a_3)\, e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3}\ da_3$

By the assumptions on $\left[f_2^3\right]_3 and\ [h^3]^3 dH_3/dz > 0$. Hence, if

$$H_3(0) = 1 - \int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3} da_3 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_3(0) < 0$ follows if $g_3^-(b_3) < [f_3^{-3}]_3(r_3)$, where $r_3$ is the solution of (12).

Indeed, since

$$R_3(s_3) = \int_0^{\infty} \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 e^{-y\frac{a}{3}}da_3}$$

Is a decreasing function of $s_3$ when $P_3^* = [g^{-3}]_3(b_3) < [f_3^{-3}]_3(r_3)$, that is $[f_3]_3(P_3^*) < r$, $R\big(f_3(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_3(g^{-3}(b))) > 1, i.e.,$

$\int_0^{\infty} \beta_3(a_3)e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3\left(g_3^{-3}(b)\right)a_3} > 1.$ (H4)

Then the system (4)-(6) has a unique positive equilibrium $(P_3^*(a_3), Q_3^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_3^*(a_3), Q_3^*)$ is more complicated. We again consider perturbations $m_3(t_3a_3) = p_3(t_3, a_3) - p_3^*(a_3)$ and $W_3(t) = Q_3(t) - Q^*$.Then

$M_3(t) \equiv \int_0^{\infty} m_3(t, a_3)da_3 = P_3(t) - P_3^*,$

$$[m_3]_3 + [m_a]_3 = [p_3]_3 + [p_a]_3 - \left[p^*\right]_3(a_3) \approx -(\mu_3(a_3) + [f_{3*}]_3 + [f_{2*}]_3)m_3$$

$$-[f_{3*}^3]_3 p_3^*(a_3)M - [f_{2*}^3]_3 p_3^*(\mathrm{a})W_3,$$

$$m_3(t,0)p_3(t,0) - p_3^*(0) \approx h_* \int_0^{\infty} \beta_3(a_3)m_3(t,a_3)\, da + \frac{[h_*^3]_3}{[h_*]_3}\, p_3^*(0)W_3$$

$$\frac{dW_3}{dt} = \frac{dQ_3}{dt} \approx [g_*^3]_3 Q_3^* M_{,3}$$

Where $[f_{3*}]_3$ denoting $[f_3(P^*)]_3$, $[f_{3*}]_3$ denoting $[f_2(P^*)]_3$, $[f'_{3*}]_3$ denoting $[f'_3(P^*)]_3[f'_{2*}]_3$ denoting $[f_2]_3$ $(Q_3^*,\quad h_3^*$ denoting $h'_*$ denoting $h'_3$ $(Q_3^*)$, and $[g'_*]_3$

Denoting $[g']_3(P_3^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_3]_3 + [m_a]_3 = -(\mu_3(a_3) + [f_{3*}]_3 + [f_{2*}]_3)m_3 - [f'_{3*}]_3[p_3^*](a_3)M_3 - [f'_{2*}]_3 p_3^*(a_3)W_3$$

$$m_3(t,0) = [h_*]_3 \int_0^\infty \beta_3(a_3)m_3(t,a_3)da_3 + \frac{[h_*^3]_3}{[h^*]_3}[p^*]_3(0)\,W_{,3} \qquad (30)$$

$$\frac{dW_3}{dt} = [g_*^3]_3[Q_3^*]M_3$$

Integrating $[m_3]_3(t,a_3)$ with respect to a from 0 to ∞, we have

$$\frac{dM_3(t)}{dt} = \int_0^\infty [mt]_3(t,a_3)da_3$$

$$= -\int_0^\infty [m_a]_3\,(t,a_3)da_3 - \int_0^\infty \mu(a_3)m(t,a_3)da_3$$

$$-\left([f_{3*}]_3 + [f_{2*}]_3 + [f'_{3*}]_3\int_0^\infty p_3^*(a_3)da_3\right)M_3 - [f_{2*}^3]_3\int_0^\infty p_3^*(a_3)da_3 W_3$$

$$= m_3(t,0) - \int_0^\infty \mu_3(a_3)\,m_3(t,a_3)da_3$$

$$-\left([f_{3*}]_3 + [f_{2*}]_3 + [f'_{3*}]_3 P_3^*\right)M_3 - [f'_{2*}]_3 P_3^* W_{,3}$$

From (30), we arrive at the equations

$$\frac{dW_3(t)}{dt} = \int_0^{\infty} ([h_*]_3\beta_3(a_3) - \mu_3(a_3))m_3(t,a_3)da_3$$

$$-\left([f_{3*}] + [f_{2*}]_3 + \left[f_{3*}^{'}\right]_3 P_3^*\right) M_3 + \left(\frac{[h_*^{'}]_3}{[h_*]^3}\, P_3^*(0) - f_{2*}^{'} P_3^*\right) W_3,$$

(33)

$$\frac{dW_3(t)}{dt}\left[g^{'}{}_*\right]_3 Q_3^*\, M_{.3}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{n*}]_3\beta_3(a_3) - \mu_3(a_3))m_3(t,a_3)da_3 = ([h_*]_3\beta_3(\theta_3) - \mu_3(\theta_3))M_{,3}$$

Where $0< \theta_3 = \theta_3(t) < \infty$.

And if $a_3^2\,[a_3^2]_3<4\,[a_2]_3<4\,a_2 e$,

$\mathrm{Re}\,\lambda_3\left[\lambda_3\right]_3 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_3 = -\frac{3}{2}a_3[a2\ ]_3 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_*\,[h_*]_3\beta_3(a_3) - \mu_3(a_3)\right|$. Then$|n_3(t)|$ ⊠⊿∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_3^2 <4\,a_2 e$ and $[a_3] > 2\Delta$,or $[a_3^2]$ ⊠$[a_2]e$ and $(a_3 - \sqrt{(a_3^2)_3 - 4(a_2)_3 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+3}\right)$ and $\beta_3(a_3) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3) + d\tau_3 da} = \frac{3}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_3 + [\mu_0]$ 0.5. Assume $[f_3](x_3) = K_3 x_3$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_3(x_3) = [0x_3]/(3 + x_3)$; the functional responses $[[f_2](x) = [0]\, x_3/(1 + x_3)$ and $h_3(x_3)=(1+\omega_3 x_3)/(1 + x_3)$ (See Holling [7] and Rashevsky [13])

Equation (34),now has the form of

$$\frac{m_3 + [\mu_0] + [K][P_0])}{[m_3]_3 + (\mu_0)_3 + [K][P_0])_3^3\ [K]_3 [P_0]_3^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(3 + P_0) < b_3$ i e $[P_0] < b/(\delta_3 - b_3)$ then the equilibrium $(P_0)$ $[P_0]_3(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_3$

$[f_{3*}] = K_3\ P^*\ [f_{2*}] = \frac{\sigma Q_3^*}{3+Q_3^*} \quad [h_*] = \frac{3+\omega_3 Q_3^*}{3+Q_3^*}$

$[f_{3*}] = K_3\ [f_{2*}^3] = \frac{\sigma_3}{(3+Q_3^*)^2} [\ g_*^3] = \frac{\sigma_3}{(3+P_3^*)^2} \quad [h_*^3] = \frac{\omega_3 - 3}{(3+Q_3^*)^2}$

In order to have (H3) and (H4) hold, $m, \mu_0, K, b, \delta$, and $\omega$ need to satisfy

$$\frac{\omega_3}{(m+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma)^2}+\frac{2\omega_3}{([m_3]_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma_3)^3}<1$$

And

$$\frac{1}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^2}+\frac{2\omega_3}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}+(m_3+[\mu_0]_3+\frac{K_3b_3}{\delta_3-b_3}+\sigma_3)\,Q_3^*)^3$$

$$=m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}(1+Q^*)(1+\omega_3Q_3^*$$

Now

$$\left[a_3^2\right]-4\left[a_2\right]e=\left(\frac{2K_3b_3}{\delta_3-b_3}\right)^2+\frac{4K_3b_3\sigma_3\,Q_3^*}{(\delta_3-b_3)(1+Q_3^*)}+\frac{\sigma_3^2Q_3^{*2}}{(1+Q_3^*)^2}$$

$$\frac{4\,b_3\sigma_3\,(\delta_3-b)Q^*}{\delta_3(1+Q_3^*)^2}-\frac{4\,(\delta_3-b_3)^2(1-\omega_3)Q_3^*}{\delta_3(1+\omega_3Q_3^*)(1+Q_3^*)}P*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_3\,ae_3^{mu}\;\;[\mu_0]_3-\frac{a_3}{a_3+3}$$

Hence if $[a_3^2 > 4(a_2)_3 e$ and

$$[a_3] = \frac{2K_3 b_3}{\delta_3 - b_3} + \frac{\sigma_3 Q_3^*}{1 + Q_3^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_3]_3 = [f_{3*}]_3 + [f_{2*}]_3 + [f_{3*}] P_3^*$$

$$[a_2]_3 = \left[f_{2*}^3\right] P_3^* - \frac{h_*^3}{h_*} P^*(0)$$

$$e = [g_*^3]_3 Q_3^*$$

$$\eta_3(t) = \left[h_*\right]_3 \beta_3 \left(\theta_3(t)\right) - \mu_3(\theta(t)),$$

The system(33)can be written as the following system:

$$[M(t)]\left[a_3\right] M_3 - [a_2]_3 W_3 - [n(t)] M_3$$

$$w_3(t) = eM_3$$

Write

$$A = \begin{pmatrix} -[a_3]_3 - [a_2]_3 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_3 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_3 \\ W_3 \end{pmatrix}$$

The solution of (32)is

$$X_3(t) = e^{A,t}[X]_3(0) + \int_0 e^{At}$$

**This equation is useful in following type of problems . For actual model kindly see various papersThere is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant Λ, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive Λ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the Λ>0 generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible.**

**Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of (e) GW signals from BBHs with component spins $\chi_{1,2} \in [-0.99, 0.99]$ aligned (e&eb) with the orbital angular momentum, component masses $m_{1,2} \in [2,48] M_\odot$, and (e&eb) total mass $M_{total} \leq 50\ M_\odot$. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins >+0.6 and component masses $\gtrsim 20\ M_\odot$. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201 STOP Advanced LIGO may be the first experiment to (e) detect gravitational waves. Through superradiance of stellar black holes, it may also be (=) the first experiment to discover the QCD axion with (e&eb) decay constant above the GUT scale. When an axion's Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e) the black hole, forming a "gravitational atom." Through the superradiance process, the number of axions occupying (e) the bound levels grows exponentially, extracting (e) energy and angular momentum from the black hole. Axions**

**transitioning between levels of the gravitational atom and (e&eb) axions annihilating to (e) gravitons can produce (eb) observable gravitational wave signals. The signals are long-lasting, monochromatic, and can be distinguished from (e) ordinary astrophysical sources. Authors estimate up to O (1) transition events at aLIGO for (e) an axion between 10^-11 and (e&eb) 10^-10 eV and up to 10^4 annihilation events for (e) an axion between 10^-13 and 10^-11 eV. In the event of a null search, aLIGO can constrain (e) the axion mass for (e) a range of rapidly spinning black hole formation rates. Axion annihilations are also promising for (e) much lighter masses at future lower-frequency gravitational wave observatories; the rates have (e) large uncertainties, dominated by (e) supermassive black hole spin distributions. Projections for aLIGO are (=) robust against perturbations from (e) the black hole environment and account for (e) updated exclusion on (e&eb) the QCD axion of 6*10^-13 eV < ma < 2*10^-11 eV suggested by (e) stellar black hole spin measurements. Subjects: High Energy Physics - Phenomenology (hep-ph); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 91, 084011 (2015) DOI: 10.1103/PhysRevD.91.084011Cite as:arXiv:1411.2263 [hep-ph] (or arXiv:1411.2263v3 [hep-ph] for this version) Discovering the QCD Axion with Black Holes and Gravitational Waves Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang STOP Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric**

**charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by (e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave**

**signal. The signal sweeps upwards in frequency from 35 to 250 Hz with (e&eb) a peak gravitational-wave strain of (e) 1.0×10−21. It matches the waveform predicted by general relativity for (e) the inspiral and merger of a pair of black holes and the ringdown of the resulting (eb) single black hole. The signal was observed with a matched-filter signal-to-noise ratio of (e) 24 and a false alarm rate estimated to be (=) less than 1 event per 203 000 years, equivalent to (=) a significance greater than 5.1σ. The source lies at (eb) a luminosity distance of 410+160−180 Mpc corresponding to (e&eb) a redshift z=0.09+0.03−0.04. In the source frame, the initial black hole masses are (=) 36+5−4M⊙ and 29+4−4M⊙, and the final black hole mass is (=) 62+4−4M⊙, with (e&eb) 3.0+0.5−0.5M⊙c2 radiated in (eb) gravitational waves. All uncertainties define (eb) 90% credible intervals. These observations demonstrate (eb) the existence of binary stellar-mass black hole systems. Observation of Gravitational Waves from a Binary Black Hole Merger B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 061102 – Published 11 February 2016 This is the first direct detection of gravitational waves and (e&eb) the first observation of a binary black hole merger. Received 21 January 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.061102 STOP This is the second in a series of three papers on (e&eb) systematic analysis of rank 1 Coulomb branch geometries of (e) four dimensional N=2 SCFTs. In the first paper authors developed a strategy for (e) classifying physical rank-1 CB geometries of N=2 SCFTs. They show how to carry out this strategy computationally to construct (eb) the Seiberg-Witten curves and (e&eb) one-forms for all the rank-1 SCFTs. Explicit expressions are given for all cases, with the exception of the Nf=4 SU (2) gauge theory and (e&eb) the En SCFTs which were previously constructed. Classification includes all known rank-1 theories plus (+) a new one with (e&eb) an abelian flavor group, plus (+) nine additional theories whose existence is more speculative. Four of those, reported in first paper, depend on (=) the assumption of new frozen rank-1 SCFTs. Here authors also show that the assumption of the existence of certain rank-0 N=2 SCFTs leads to (eb) five additional consistent**

**rank-1 CB geometries. Subjects: High Energy Physics - Theory (hep-th) Cite as:arXiv: 1601.00011 [hep-th] (or arXiv: 1601.00011v2 [hep-th] for this version Geometric constraints on the space of N=2 SCFTs II: Construction of special Kähler geometries and RG flows Philip C. Argyres, Matteo Lotito, Yongchao Lü, Mario Martone (Submitted on 31 Dec 2015 (v1), last revised 6 Jan 2016 (this version, v2)) STOP Authors propose a model of Dark Energy in which the field currently dominating (e) the energy density of the universe is (=) an "axion field" linearly coupled to (e&eb) the Pontryagin density, tr (FΛF), (i.e., the exterior derivative of (e) the Chern-Simons form) of (e) a massive gauge field. Authors assume that the axion has (e) self-interactions corresponding to (e&eb) a non-trivial (exponential) potential. Authors argue that a non-vanishing magnetic helicity of (e) the gauge field triggers (eb) slow-rolling of the axion at field values far below (eb) the Planck scale. Proposal leads to (eb) a "Tracking Dark Energy Scenario" in which the contribution of the axion energy density to (e) the total energy density is (=) constant (and small) during (e&eb) the early radiation phase, until (e) a secular growth term proportional to (e) the Pontryagin density of (e) the gauge field becomes (=) dominant. The initially small contribution of (e) the axion field to (e) the total energy density is related to the observed small baryon-to-entropy ratio. Some speculations concerning the nature of the gauge field are offered.**

**Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Cite as:arXiv: 1601.00057 [hep-th] (or arXiv:1601.00057v2 [hep-th] for this version) Tracking Dark Energy from Axion-Gauge Field Couplings Stephon Alexander, Robert Brandenberger, Juerg Froehlich (Submitted on 1 Jan 2016 (v1), last revised 23 Jan 2016 (this version, v2)) STOP Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a**

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**associated entropic measure. Subjects: High Energy Physics - Theory (hep-th) DOI:10.1016/j.physletb.2016.02.038 Cite as: arXiv: 1601.00076 [hep-th] (or arXiv: 1601.00076v2 [hep-th] for this version) Bounds on topological Abelian string-vortex and string-cigar from information-entropic measure R. A. C. Correa, D. M. Dantas, C. A. S. Almeida, Roldao da Rocha (Submitted on 1 Jan 2016 (v1), last revised 17 Feb 2016 (this version, v2)) STOP Sum of the Holevo quantity (that bounds the capacity of quantum channels to transmit (e&eb) classical information about an observable) and (e&eb) the quantum discord (a measure of the quantumness of (e) correlations of that observable) yields (eb) an observable-independent total given by (e) the quantum mutual information This split naturally delineates ,dilates and disseminates information about (e&eb) quantum systems accessible to (e) observers -- information that is redundantly transmitted by (e&eb) the environment -- while showing that it is (=) maximized for (e) the quasi-classical pointer observable. Other observables are accessible only via (e&eb) correlations with the pointer observable .Further authors prove an anti-symmetry property relating accessible information and (e&eb) discord. It shows that information becomes (=) objective -- accessible to (e) many observers -- only as quantum information is relegated to (e) correlations with (e&eb) the global environment, (eb) and, therefore, locally inaccessible. The resulting complementarity explains (eb) why, in a quantum Universe, we perceive objective classical reality while (e) flagrantly quantum superpositions are out of reach. Subjects: Quantum Physics (quant-ph) Journal reference: Scientific Reports 3, 1729 (2013) Report number: LAUR-13-26386 Cite as:arXiv: 1303.4659 [quant-ph] (or arXiv: 1303.4659v2 [quant-ph] for this version) Complementarity of quantum discord and classically accessible information Michael Zwolak, Wojciech H. Zurek STOP Niels Bohr proposed that the outcome of the measurement becomes (=) objective and real, and, hence, classical, when (e) its results can be communicated by (e) classical means. In this work authors revisit Bohr's postulate using (e) modern tools from (e) the quantum information theory. They find a**

**full confirmation of Bohr's idea: if a measurement device is in a nonclassical state, (eb) the measurement results cannot be communicated perfectly by classical means. In this case some part of information in the measurement apparatus is lost (e) in the process of communication: the amount of this lost information turns out to be (=) the quantum discord. The information loss occurs even when (e) the apparatus is not entangled with the system of interest. The tools presented in this work allow (eb) to generalize Bohr's postulate: authors show that for pure system-apparatus states (eb) quantum communication does not provide any advantage when (e) measurement results are communicated to (e) more than one recipient. They further demonstrate the superiority of quantum communication to (e&eb) two recipients on (e&eb) a mixed system-apparatus state and show that (eb) this effect is fundamentally different from (e) quantum state cloning. Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. Lett. 111, 040401 (2013) DOI:10.1103/PhysRevLett.111.040401 Report number: LA-UR-12-25351 Cite as: arXiv: 1304.2914 [quant-ph] (or arXiv: 1304.2914v2 [quant-ph] for this version) Quantum discord cannot be shared Alexander Streltsov, Wojciech H. Zurek STOP Wojciech H. Zurek, Don N. Page investigate thermodynamic equilibrium of (e) a self-gravitating perfect fluid in (eb) a spherically symmetric system containing (e) a black hole of mass M by means of (e) the Tolman-Oppenheimer-Volkoff (TOV) equation. At r >> 2M its solutions describe (eb) a black-body radiation atmosphere with (e&eb) the Hawking temperature T_BH~1/(8 \pi M) that is increasingly blueshifted as r approaches (eb) 2M. However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments:**

**Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by**

**(e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a firewall-like "Planck cocoon" with entropy close to (e) black hole entropy**

# Prescientiаl Dynamics And Contential Staticity Of Accentuation Attrition Models

## Module Four

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equailibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-445) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Some of the concepts to which model is applicable is as follows:**

**Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by (e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave signal. The signal sweeps upwards in frequency from 35 to 250 Hz with (e&eb) a peak gravitational-wave strain of (e) 1.0×10−21. It matches the waveform predicted by general relativity for (e) the inspiral and merger of a pair of black holes and the ringdown of the resulting (eb) single black hole. The signal was observed with a**

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Almeida, Roldao da Rocha (Submitted on 1 Jan 2016 (v1), last revised 17 Feb 2016 (this version, v2)) STOP Sum of the Holevo quantity (that bounds the capacity of quantum channels to transmit (e&eb) classical information about an observable) and (e&eb) the quantum discord (a measure of the quantumness of (e) correlations of that observable) yields (eb) an observable-independent total given by (e) the quantum mutual information This split naturally delineates ,dilates and disseminates information about (e&eb) quantum systems accessible to (e) observers -- information that is redundantly transmitted by (e&eb) the environment -- while showing that it is (=) maximized for (e) the quasi-classical pointer observable. Other observables are accessible only via (e&eb) correlations with the pointer observable .Further authors prove an anti-symmetry property relating accessible information and (e&eb) discord. 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However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments: Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by (e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a**

**firewall-like "Planck cocoon" with entropy close to (e) black hole entropy. Authors reproduce paper with an updated title and abstract. For a later study with similar results see G. 't Hooft, gr-qc/9706058 Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Physical Review D, Volume 29, Number 4, pp. 628 - 631 (1984) DOI:10.1103/PhysRevD.29.628 Cite as:arXiv:1511.07051 [gr-qc] (or arXiv:1511.07051v1 [gr-qc] for this version) Firewalls, black-hole thermodynamics, and singular solutions of the Tolman-Oppenheimer-Volkoff equation Wojciech H. Zurek, Don N. Page (Submitted on 22 Nov 2015) STOP Frederic Magniez, Ashwin Nayak consider the problem of testing the commutativity of (e) a black-box group specified by its k generators. The complexity (in terms of k) of (e) this problem was first considered by Pak, who gave (eb) a randomized algorithm involving O(k) group operations. They construct a quite optimal quantum algorithm for this problem whose complexity is in O~(k2/3)O~(k2/3). The algorithm uses and highlights the power of the quantization method of Szegedy. For the lower bound of Ω(k2/3)Ω(k2/3), authors give a reduction from (e) a special case of Element Distinctness to our problem. Along the way, they prove the optimality of the algorithm of (e) Pak for the randomized model. Algorithmica July 2007, Volume 48, Issue 3, pp 221-232 First online: 25 June 2007 Quantum Complexity of Testing Group Commutativity Frederic Magniez, Ashwin Nayak STOP Pawel Wocjan, Jon Yard analyze relationships between quantum computation and (e&eb) a family of generalizations of the Jones polynomial. Extending recent work by Aharonov et al., authors give efficient quantum circuits for (e) implementing the unitary Jones-Wenzl representations of the braid group. We use these to provide new quantum algorithms for (e) approximately evaluating a family of specializations of the HOMFLYPT two-variable polynomial of trace closures of braids. They also give (eb) algorithms for approximating the Jones polynomial of a general class of closures of (e) braids at roots of unity. Next authors provide (eb) a self-contained proof of a result of (e) Freedman et al. that any quantum computation can be replaced by (e&eb) an additive approximation of the Jones polynomial, evaluated at (eb) almost any primitive root of unity. Proof encodes two-qubit unitaries into (e&eb) the rectangular representation of (e) the eight-strand braid group. Authors then give QCMA-complete and (e&eb) PSPACE-complete problems which are based on braids. They conclude with direct proofs that evaluating the Jones polynomial of the plat closure at most primitive roots of unity is a #P-hard problem, while learning its most significant bit is PP-hard, circumventing (e) the usual route through the Tutte polynomial and graph coloring. Subjects: Quantum Physics (quant-ph); Geometric Topology (math.GT) Cite as: arXiv:quant-ph/0603069 (or arXiv:quant-ph/0603069v3 for this version) The Jones polynomial: quantum algorithms and applications in quantum complexity theory Pawel Wocjan, Jon Yard STOP Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan address the task of switching off the Hamiltonian of (e) a system by removing (e) all internal and system-environment couplings. Authors propose dynamical decoupling schemes, which use (e) only bounded-strength controls, for quantum many-body systems with local system Hamiltonians and local environmental couplings. To do so, they introduce (eb) the combinatorial concept of (e) balanced-cycle orthogonal arrays (BOAs) and show how to construct them from (e) classical error-correcting codes. The derived decoupling schemes may be useful as (=) a primitive for more complex schemes, e.g., for Hamiltonian simulation. For the case of n qubits and a 2-local Hamiltonian, the length of the resulting decoupling (e&eb) scheme scales as O (nlogn), improving over (e&eb) the previously best-known schemes that scaled quadratically with (e&eb) n. More generally, using balanced-cycle orthogonal arrays constructed from families of BCH codes, we show that bounded-strength decoupling for any ℓ-local Hamiltonian, where ℓ≥2, can be achieved using decoupling schemes of length at most O(nℓ−1logn). Subjects:Quantum Physics (quant-ph); Emerging Technologies (cs.ET) Journal reference: IEEE Trans. Inf. Theory, 62(5): 2881-2894, 2016 DOI: 10.1109/TIT.2016.2535183 Report number: MIT-CTP /4705 Cite as: arXiv: 1509.00408 [quant-ph] (or arXiv: 1509.00408v2 [quant-ph] for this version) Improved bounded-strength decoupling schemes for local Hamiltonians Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan STOP Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola propose dynamical control schemes for Hamiltonian simulation in many-body quantum systems that avoid instantaneous control operations and rely solely on (e) realistic bounded-strength control Hamiltonians. Each simulation protocol consists of (e) periodic repetitions of a basic control block, constructed as (=) a suitable modification of an "Eulerian decoupling cycle," that would otherwise implement (eb) a trivial (zero) target Hamiltonian. For an open quantum system coupled to (e) an uncontrollable environment, approach may be employed to (e) engineer an effective evolution that simulates (e&eb) a target Hamiltonian on the system, while suppressing (e) unwanted decoherence to the leading order. Authors present illustrative applications to (e) both closed- and open-system simulation settings, with (e&eb)**

**emphasis on simulation of non-local (two-body) Hamiltonians using (e) only local (one-body) controls. In particular, authors provide (eb) simulation schemes applicable to (e&eb) Heisenberg-coupled spin chains exposed to (e) general linear decoherence, and show how to simulate (e&eb) Kitaev's honeycomb lattice Hamiltonian starting from Ising-coupled qubits, as (=) potentially relevant to (e&eb) the dynamical generation of a topologically protected (e&eb) quantum memory. Additional implications for quantum information processing are discussed. Subjects: Quantum Physics (quant-ph) Journal reference: New J. Phys. 16, 045021 (2014) DOI: 10.1088/1367-2630/16/4/045021 Cite as:arXiv: 1310.4153 [quant-ph]**

**(or arXiv: 1310.4153v1 [quant-ph] for this version) Hamiltonian quantum simulation with bounded-strength controls Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola STOP Felix Fontein, Pawel Wocjan present a quantum algorithm for computing the period lattice of (e) infrastructures of fixed dimension. The algorithm applies to (e&eb) infrastructures that satisfy certain conditions. The latter are always fulfilled for (e) infrastructures obtained from (e) global fields, i.e., algebraic number fields and function fields with (e&eb) finite constant fields. The first main contributions is an exponentially better method for (e) sampling approximations of vectors of the dual lattice of (e) the period lattice than the methods outlined in the works of (e) Hallgren and Schmidt and Vollmer. This new method improves (eb) the success probability by (e) a factor of at least 2^ {n^2-1} where (e) n is the dimension. The second main contribution is a rigorous and complete proof that the running time of the algorithm is (=) polynomial in the logarithm of (e) the determinant of the period lattice and (e&eb) exponential in n. The third contribution is the determination of (e) an explicit lower bound on (e&eb) the success probability of our algorithm which greatly improves on (e&eb) the bounds given in the above works. The exponential scaling seems inevitable because (e) the best currently known methods for (e) carrying out fundamental arithmetic operations in infrastructures obtained from (e) algebraic number fields take exponential time. In contrast, the problem of computing (e&eb) the period lattice of infrastructures arising from function fields can be solved without (e) the exponential dependence on (e&eb) the dimension n since this problem reduces efficiently to the abelian hidden subgroup problem. This is also true for (e) other important computational problems in (eb) algebraic geometry. The running time of the best classical algorithms for (e) infrastructures arising from (e) global fields increases subexponentially with (e&eb) the determinant of the period lattice. Subjects: Quantum Physics (quant-ph) Cite as: arXiv: 1111.1348 [quant-ph]**

**(or arXiv: 1111.1348v2 [quant-ph] for this version) Quantum Algorithm for Computing the Period Lattice of an Infrastructure Felix Fontein, Pawel Wocjan STOP Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan advocate a new approach of (e) addressing hidden structure problems and (e&eb) finding efficient quantum algorithms. Authors introduce and investigate the Hidden Symmetry Subgroup Problem (HSSP), which is a generalization of the well-studied Hidden Subgroup Problem (HSP). Given a group acting on (e&eb) a set and an oracle whose level sets define (eb) a partition of the set; the task is to recover (e) the subgroup of symmetries of (e) this partition inside the group. The HSSP provides (eb) a unifying framework that, besides the HSP, encompasses (e) a wide range of algebraic oracle problems, including (e) quadratic hidden polynomial problems. While the HSSP can have (e) provably exponential quantum query complexity, authors obtain efficient quantum algorithms for various interesting cases. To achieve this, authors present a general method for (e) reducing the HSSP to the HSP, which works efficiently in several cases related to (e&eb) symmetries of polynomials. The HSSP therefore connects in a rather surprising way certain hidden polynomial problems with (e&eb) the HSP. Using this connection, authors obtain the first efficient quantum algorithm for (e) the hidden polynomial problem for (e) multivariate quadratic polynomials over fields of (e) constant characteristic. Authors also apply the new methods to (e&eb) polynomial function graph problems and (e&eb) present an efficient quantum procedure for (e) constant degree multivariate polynomials over (e&eb) any field. This result improves in several ways the currently known algorithms. Subjects: Quantum Physics (quant-ph) Journal reference: SIAM Journal on Computing 42:5 (2013) DOI: 10.1137/120864416 Cite as: arXiv: 1107.2189 [quant-ph]**

**(or arXiv: 1107.2189v2 [quant-ph] for this version) Hidden Symmetry Subgroup Problems Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan STOP Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan present an efficient general method for realizing a quantum walk operator corresponding to (e&eb) an arbitrary sparse classical random walk. Approach is based on Grover and Rudolph's method for preparing coherent versions of (e) efficiently integrable probability distributions. This method is intended for (e) use in quantum walk algorithms with polynomial speedups, whose complexity is usually measured in terms of (e&eb) how many times we have to apply a step of a quantum walk, compared to (e&eb) the number of necessary classical Markov chain steps. Authors consider a finer notion of complexity including (e) the**

**number of elementary gates it takes to implement each step of the quantum walk with (e&eb) some desired accuracy. The difference in complexity for (e) various implementation approaches is that method scales linearly in the sparsity parameter and (e&eb) poly-logarithmically with (e&eb) the inverse of the desired precision. Best previously known general methods either scale quadratically in (eb) the sparsity parameter, or polynomially in (eb) the inverse precision. Approach is especially relevant for (e) implementing quantum walks corresponding to (e&eb) classical random walks like those used in (eb) the classical algorithms for approximating permanents and sampling from binary contingency tables. In those algorithms, the sparsity parameter grows with (e&eb) the problem size, while (e) maintaining high precision is required. Subjects: Quantum Physics (quant-ph) Journal reference: Quantum Information and Computation, Vol.10, No.5&6, pp0420-0434 (2010) Cite as: arXiv: 0903.3465 [quant-ph] (or arXiv: 0903.3465v3 [quant-ph] for this version) Efficient Circuits for Quantum Walks Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan STOP Preparing the ground state of a system of interacting classical particles is (=) an NP-hard problem. Thus, there is in general no better algorithm to (e) solve this problem than exhaustively going through all N configurations of (e) the system to determine (eb) the one with lowest energy, requiring (e) a running time proportional to N. A quantum computer, if it could be built, could solve (e) this problem in time sqrt (N). Here, authors present a powerful extension of this result to (e) the case of interacting quantum particles, demonstrating (eb)that a quantum computer can prepare the ground state of (e) a quantum system as efficiently as it does for (e) classical systems. Subjects: Quantum Physics (quant-ph) Journal reference: Phys.Rev.Lett.102:130503, 2009 DOI:10.1103/PhysRevLett.102.130503 Cite as: arXiv: 0809.2705 [quant-ph] (or arXiv: 0809.2705v1 [quant-ph] for this version) Preparing ground states of quantum many-body systems on a quantum computer David Poulin, Pawel Wocjan STOP Daniel Nagaj, Pawel Wocjan construct a simple translationally invariant, nearest-neighbor Hamiltonian on (e&eb) a chain of 10-dimensional qudits that makes it possible to (e) realize universal quantum computing without (e) any external control during (e&eb) the computational process. Authors only require the ability to prepare an initial computational basis state which encodes (e) both the quantum circuit and its input. The computational process is (=) then carried out by (e) the autonomous Hamiltonian time evolution. After a time polynomially long in the size of (e) the quantum circuit has passed, the result of (e) the computation is obtained with (e&eb) high probability by (e) measuring a few qudits in the computational basis. This result also implies (eb) that there cannot (e) exist efficient classical simulation methods for (e) generic translationally invariant nearest-neighbor Hamiltonians on (e&eb) qudit chains, unless (e) quantum computers can be efficiently simulated by (e&eb) classical computers (or, put in complexity theoretic terms, unless BPP=BQP). Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. A 78, 032311 (2008) DOI: 10.1103/PhysRevA.78.032311 Cite as: arXiv: 0802.0886 [quant-ph] (or arXiv: 0802.0886v4 [quant-ph] for this version) Hamiltonian Quantum Cellular Automata in 1D Daniel Nagaj, Pawel Wocjan**

## 28. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_4}{dt} = -b_4 Q_4 + g_4(g_4)Q_4$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_4(0) = 0, \quad g_4(P_4) > 0, \; P_4 \; 0, \qquad \lim_{P_4 \to \infty} g_4(P_4) = (K_4)_4 \; \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_4 + (\rho_a)_4 + D_4(a_4, P_4, Q_4)\rho_4 + 0, \qquad 0 < a_4 \; \infty,$$

$$\rho_4(t, 0) = \int_0^\infty n_4(a_4, P_4, Q_4)\rho_4(t_4, a_4)da_4$$

Where $D_4$ is the mortality function and $n_4$ is the fertility function.

Put $_4$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_4}{dt}\right]_4 = x_4\left[\frac{dQ_2}{dt}\right]_4 - x_2\left[\frac{dQ_4}{dt}\right]_4$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_4 - x_4\left[\frac{dQ_2}{dt}\right]_4$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_4 - x_6\left[\frac{dQ_3}{dt}\right]_4$$

Where $x_4, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_4$ and $g_4(P_4)Q_4$

Now when you add the term $(b_4Q_4)$ with $g_4(P_4)(Q_4)$ we are infect see is that $(b_4Q_4)$ is credited incrementally part of $g_4(P_4)(Q_4)$ or instant nearly, when subtracted from $\frac{dQ_4}{dt}$, it means that $g_4(P_4)(Q_4)$ is slowly eaten by $\frac{dQ_4}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_4 = \mu_4(a_4) + (f_4)_4(P_4)_4 + (f_2)_4(Q)_4 \quad (2)$$

$$n_4 = \beta_4(a_4)h_4(Q_4), \quad (3)$$

Where $\mu_4$the age- dependent debit rate whit is $\mu_4(a_4)$ $\forall a > 0$; $(f_4)_4$ is the density- dependent death rate with $f_4(0) = 0$, $(f_4^4)_4(P_4)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_4)_4(P_4) = (k_4)_4$ ; $f_2$ is the functional response of creditors to prey with $[f_2]_4(0) = 0, (f_2^4)_4(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_4(Q_4) = [k_2]$ ; $\beta$ is the age dependent birth rate with $\beta(a)$ $0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 4, h^4(Q)$ $0, \forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_4 + (\rho_a)_4 = -\mu_4(a_4) + (f_4)_4P_4) + \big(f_2)_4(Q_4)\big)(\rho_4)_4 \quad (4)$$

$$\rho_4(t_4, 0) = h_4(Q_4)\int_0^\infty \beta_4(a_4)\rho_4(t, a_4)da, \quad (5)$$

$$\frac{dQ_4}{dt}$$

$$= -b_4Q_4 + g_4(P_4)Q_4 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_4$ and $Q_4$ of (4-6) satisfy

$$0 \leq \rho_4(t, a_4) \leq \rho_4^*(t, a_4),$$

$$0 \leq Q_4(t) \leq Q_4^*(t),$$

Where $\rho_4^*(t, a_4)$ and $Q_4^*(t)$, are solutions of

$$(\rho_t^*)_4 + (\rho_a^*)_4 = -(\mu_4(a_4) + (f_4)_4(P_4^*) + (f_2)_4 Q_4^*))\, \rho_4^*, \quad (7)$$

$$\rho_4^*(t, 0) = \int_0^\infty \beta_4(a_4)\, \rho_4^*(t, a_4) da_4 \quad (8)$$

$$\frac{dQ_4^*}{dt} = -b_4 Q_4^* + g_4(P_4^*)Q_4^*, \quad (9)$$

With $P_4^*(t) = \int_0^\infty \rho_4^*(t, a_4) da_4$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_4^*(t)}{dt} = P_4^*(t)\big(r_4 - (f_4)(P_4^*) - (f_2)_4(Q^*)\big), \quad (10)$$

$$\frac{dQ_4^*(t)}{dt} = Q_4^*(t)(-b_4 + g_4(P_4^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_4(a) e^{-\int_0^a \mu_4(\tau)dt} e^{-ra}\, da_4 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_4^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^{\infty} \beta_4(a_4) e^{-\int_0^a \mu_4(t)dt}\, da_4 > 1, \qquad (H1)$$

Which guarantees $r_4 > 0$.

*Remark.* If $\beta_4(a) = [\beta_0]_4 a_4 e^{-\alpha_4 a_4}, \mu_4(a_4) \equiv \mu_4$ then the solution of (12) is

$$r_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_4]_4(P_4) = 0, [f_2](Q_4) = \lambda_4 Q_4\ h_4(Q_4) = 1, \qquad and\ \ g_4(P_4) = c_4 P_4$$

$$[\rho_4]_4 + [\rho_a]_4 = -(\mu_4 + \lambda_4 Q_4)\rho_4$$

$$\rho_4(t,0) = \int_0^{\infty} [\beta_0]_4 a_4 e^{-\alpha_4 a_4} \rho_4(t, a_4)\, da_4$$

$$\frac{dQ_4}{dt} = -b_4 Q_4 + c_4 P_4 Q_4$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_4}{dt} = \omega_4 P_4 - \lambda_4 P_4 Q_4$$

$$\frac{dQ_4}{dt} = b_4 Q_4 - c_4 P_4 Q_4$$

Where $\omega_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$

## 29. THE EQUILIBRIUM $(\rho_0)_4(a_4), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_4(a_4), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_4 + (\rho_a)_4 = -\big(\mu_4(a_4) + (f_4)_4(P_4)\big)\rho,$$

$$\rho_4(t,0) = \int_0^\infty \beta_4(a_4)\rho_4(t,a_4)da_4 \qquad (13)$$

The solution of (13) has the form of

$$\rho_4(t,a_4) = Ce_4^{r(t-a_4)} - \int_0^a (\tau)d\tau - \int_0^4 (f_4)_4\big(P_4(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_4(t) = C_4 e^{r_4 t - \int_0^t (f_4)_4(P_4(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_4^{-4}(r)$ if $k_4 > r$. This is equivalent to

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{(k_4)_4 a_4} da_4 < 1.$$

In fact, if we define

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau)d\tau} e^{-s_4 a_4}\, da_4$$

$R(s)$ is a decreasing function. $R_4(r_4) = 1$ And $R_4(k_4)_4 < 1$ imply $r_4 < (k_4)_4$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_4(a_4), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{k_4 a_4} da_4$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_4 \equiv \int_0^\infty (\rho_0)_{44}(a_4)da_4$ *satisfies*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{-(f_4)_4(P_0)a_4} da$$

$$= 1 \qquad (11)$$

Next, we study the stability of $(\rho_0)_4(a_4), 0)$

Let $u_4(t, a_4) = \rho(t, a_4) - \rho_0(a_4)$ Then

$$N_4(t) \equiv \int_0^\infty u_4(t, a_4)da_4 = \int_0^\infty \rho_4(t, a_4)da_4 - \int_0^\infty (\rho_0)_4(a_4)da = P_4(t) - P_0)_4$$

And

$$(u_4)_4 + (u_a)_4 = (\rho_t)_4 + (\rho_a)_4 - (\rho_0^4)_4(a_4)$$

$$= -\mu_4(a_4) + (f_4)_4(P_4) + (f_2)_4(Q_4))\,(u_4 + (\rho_0)_4) + \mu_4(a_4)(\rho_0)_4 + (f_4)_4(P_0)_4(\rho_0)_4$$

$$= -\mu_4(a_4)u_4 - (f_4)_4(P_4) + (f_2)_4(Q_4))\,u_4 - (\rho_0)_4\,(f_4)_4(P_4) - (f_4)_4(P_0)_4$$

$$+ (f_2)_4(Q_4)), \qquad (15)$$

Where $(\rho_0^4)_4(a_4)$ denotes $d(\rho_0)_4(a_4)/da_4$

Expansion $(f_4)_4(P_4)$ and $(f_2)_4(Q_4)$ at $(P_0)_4$ and 0 yield respectively

$$(f_4)_4(P_4) = (f_4)_4((P_0)_4 + N_4)$$

$$= (f_4)_4(P_0)_4 + (f_4^4)_4(P_0)_4 N_4 + o_4(N_4), \quad (16)$$

$$(f_2)_4(Q_4)$$

$$= (f_2^4)_4(0)Q + o_4(Q_4) \quad (17)$$

Write $f_4(P_0) = f_{40}, f_4'(P_0) = f_{40}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_4)_4 + (u_a)_4 \approx -(\mu_4(a_4) + (f_{40})_4)u - (f_{40}^4)_4(\rho_0)_4(N_4) - (f_{20}^4)_4(\rho_0)_4(Q_4)$$

$$u_4(t,0) = \rho_4(t,0)(\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a) + (\rho_0)_4(a_4)da_4 - (\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a_4)da_4 + (h_4(Q_4) - 4)(\rho_0)_4(0) \quad (18)$$

Expanding

$$h_4(Q_4) = 1 + (h^4)_4(0)Q_4 + 0(Q_4) \equiv 1 + h_0'Q_4 + 0Q_4$$

And substituting it into (48), we arrive at

$$u_4(t,0) \approx \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + (h_{a_4}')(\rho_0)_4 Q_4$$

The linearization is the following :

$$[u_4] + [u_4] = -(\mu_4(a_4) + [t_{40}]_4 u_4 - [\alpha_4]_4(a_4)N_4 - [\alpha_2]_4(a_4)Q_4$$

$$u_4(t,0) = \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + [\xi_4]Q_4 \quad (19)$$

$$\frac{dQ_4}{dt} = -b_4 Q_4$$

where

$$[\alpha_4]_4(a_4) = [t_4^4 a]_4[p, o]_4(a_4) d2(a_4) = [f_{20}^4]_4(\rho\, o\,)_4(a_4)\xi = [\mathrm{h\ o}^4]_4[\rho \mathrm{o}]_4(\mathrm{o}) \text{ and } \mathrm{b}_4 = \mathrm{b}_4 - g_4(P_4)$$

If $g_4(P_0) > b$, clearly, $(\rho_0)_4(a_4), 0)$ is unstable and the predator population is persistent. Assume $g_4(P_0)_4 < b$.

Then, $Q_4(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_4 + u_a = -(\mu_4(a_4) + (f_{40})_4)u_4 - (\alpha_4)_4(a_4)N_4$$

$$= -(\mu_4(a_4) + f_{40})u_4 - \alpha_4(a_4)\int_0^\infty u_4(t, a_4)da_4$$

$$u_4(t, 0) = \int_0^\infty \beta_4\,(a_4)u_4(t, a_4)da_4$$

Using the method developed by Cushing [2, 3], we suppose

$$u_4(t, a_4) = \tilde{\mathrm{u}}_4(a_4)^{e^{\gamma(t-a_4)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_4}\, da_4$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_4)$ gives,

$$\tilde{\mathrm{u}}_4(a_4) = e^{-\int_0^a \mu_4(\tau_4)d\tau - (f_{40})_4 a}\left(\tilde{\mathrm{u}}_4(0) - w\int_0^a e^{\int_0^\tau \mu_4(s_4)ds + (f_{40})}\tau_4\,(\alpha_4)_4(\tau_4)e^{\gamma_4\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1-\int_0^{\infty}\beta_4(a_4)e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\right)\tilde{u}_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a(\alpha_4)_4(\tau_4)e^{-\gamma_4\tau_4}e^{(f_{40})_4\tau_4+\int_0^\tau \mu_4(s_4)ds}d\tau_4\,da_4$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}da\right)(\tilde{u}_4(0))_4$$

$$+\left(-4-\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a e^{(f_{40})_4\tau_4+\int_0^\tau \mu_4(s_4)ds+y\tau}\alpha_4(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(t_4)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_4\,(a_4)Y_4(\gamma_4,a_4)da_4\right)u_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)\int_0^a(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4=0,$$

(24)

$$\left(\int_0^{\infty}Y_4(\gamma_4,a_4)da_4\right)\tilde{u}_4(0)$$

$$+\left(-1-\int_0^{\infty} Y_4(\gamma_4,a_4)\int_0^{a}(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4\right)w=0.$$

Equqtion (20) has a solution of the form of $u_4(t,a_4)=u_4(a_4)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_4(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_4(\gamma_4,a_4)=\int_0^a(\rho_0)_4(\tau_4)Y^{-4}(\gamma_4,a_4)d\tau_4$. We obtain the characteristic equation

$$f_{40}^{l}{}_4\left(\int_0^{\infty} Y_4(\gamma_4,a_4)da_4\right)\left(\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$+\left(1-\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)da\right)\left(1+(f_{40}^{l})_4\int_0^{\infty}Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_4(a_4)=(\rho_0)_4(0)e^{-\int_0^a\mu_4(\tau_4)d\tau_4-(f_{40})_4a_4}$$

$$I_4(\gamma_4,a_4)=\int_0^a(\rho_0)_4(\tau_4)e^{\gamma_4\tau_4+(f_{40})_4\tau+\int_0^{\tau}\mu_4(s_4)ds_4}d\tau$$

$$=\frac{(\rho_0)_4(0)}{\gamma_4}(e^{\gamma_4a_4}-1)$$

Let

$$E_4(a_4)=e^{-(f_{40})_4a_4-\int_0^a\mu_4(\tau_4)d\tau_4}$$

Then $[Y_4]_4[\gamma_4]_4, a_4) = E_4(a_4)e^{-\gamma_4 a_4}$. by a simple calculation,

$$\int_0^\infty \beta_4(a_4)Y_4(\gamma_4, a_4)I(\gamma_4, a_4)da = \frac{[\rho_0]_4(0)}{\gamma}\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}\ da\right);$$

$$\int_0^\infty Y_4(\gamma_4, a_4)I(\gamma_4, a_4)da = \frac{[\rho_0]_4(0)}{\gamma_4}\left(\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}\ da_4\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_4(0)}{\gamma_4}[f_{40}^4]\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}da_4\ (\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4)$$

$$+\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right).$$

$$\left(1 + \frac{(\rho_0)_4(0)}{\gamma_4}(f_{40}^{'})_4\left(\int_0^\infty E_4(\ a_4)da_4 - \int_0^\infty E_4(\ a_4)e^{-\gamma_4 a_4}da_4\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right)(1 + \frac{(f_{40}^{'})_4}{\gamma_4}(\rho_0)_4(0)\int_0^\infty E_4(\ a_4)da_4) = 0.$$

We need ũ(a) $\not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1 - \int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)da_4 = 1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4 \neq 0,$$

Since $\int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)\ \int_0^a \alpha_4\ (\tau_4)Y^{-4}(\gamma_4, \tau_4)d\tau_4 da_4 > 0$. Therefore, the characteristic solution is

$$\gamma_4 = -(f_{40}^{'})_4(\rho_0)_4 \int_0^{\infty} E_4(a_4)da_4.$$

Since $f_4'(P_0)_4 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_4), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_4), 0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_4 + (\mu_4(a_4) + [f_4]_4(P_4^*) + [f_2]_4(Q_4^*))P_4^* = 0, \quad (26)$$

$$p_4^*(0) = h_4(Q_4^*)\int_0^{\infty} \beta_4(a_4)p_4^*(a_J)p_4^*(a_4)da_4 \quad (27)$$

$$g_4(P_4^*) - b_4 = 0, \quad (28)$$

Where $P_4^* = \int_0^{\infty} p_4^*(a_4)\, da_4$

Equation (28) yields $P^* = g^{-4}(b)$,

From (26),

$$p_4^*(a) = p_4^*(0)e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}$$

Substituting into (27), we need to solve

$$\left(1 - h_4(Q_4^*)\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}da_4\right)p_4^*(0) = 0$$

For $Q_4^*$

Fix $p_4^*(0)$ and $P_4^*$ Let

$$H_4(z_4) = 1 - h_4(z_4)\int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}da_{,4}$$

And assume that

$$[k_3]_4\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(g_4^{-4}(b_4))a_4 - [k_2]_4 a_4}\, da_4 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_4(z_4)] > 0.$$

Clearly,

$$\frac{dH_4}{dz_4} = h_4(z)(\,[f_2']_4(z_4)\,)\int_0^\infty a_4\beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - f_4(P_4^*)a_4 - f_2(z_4)a_4}da_{\square}$$

$-h'(z_4)\int_0^\infty \beta_4(a_4)\, e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}\quad da_4$

the assumptions on $\left[f_2^4\right]_4 and\ [h^4]^4 dH_4/dz > 0$. Hence, if

$$H_4(0) = 1 - \int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4} da_4 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_4(0) < 0$ follows if $g_4^-(b_4) < [f_4^{-4}]_4(r_4)$, where $r_4$ is the solution of (12).

Indeed, since

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 e^{-y\frac{a}{4}} da_4}$$

Is a decreasing function of $s_4$ when $P_4^* = [g^{-4}]_4(b_4) < [f_4^{-4}]_4(r_4)$, that is $[f_4]_4(P_4^*) < r$, $R\left(f_4(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_4(g^{-4}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4\left(g_4^{-4}(b)\right)a_4} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_4^*(a_4), Q_4^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_4^*(a_4), Q_4^*)$ is more complicated. We again consider perturbations $m_4(t_4 a_4) = p_4(t_4, a_4) - p_4^*(a_4)$ and $W_4(t) = Q_4(t) - Q^*$.Then

$$M_4(t) \equiv \int_0^\infty m_4(t, a_4) da_4 = P_4(t) - P_4^*,$$

$$[m_4]_4 + [m_a]_4 = [p_4]_4 + [p_a]_4 - [p^*]_4(a_4) \approx -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4$$

$$-[f_{4*}^4]_4 p_4^*(a_4)M - [f_{2*}^4]_4 p_4^*(\mathrm{a})W_4,$$

$$m_4(t,0)p_4(t,0) - p_4^*(0) \approx h_* \int_0^\infty \beta_4(a_4)m_4(t,a_4)\,da + \frac{[h_*^4]_4}{[h_*]_4}\, p_4^*(0)W_4$$

$$\frac{dW_4}{dt} = \frac{dQ_4}{dt} \approx [g_*^4]_4 Q_4^* M_{,4}$$

Where $[f_{4*}]_4$ denoting $[f_4(P^*)]_4$, $[f_{4*}]_4$ denoting $[f_2(P^*)]_4$, $[f_{4*}^{'}]_4$ denoting $[f^{'}{}_4(P^*)]_4 [f_{2*}^{'}]_4$ denoting $[f_2]_4$ $(Q_4^*,$ $h_{*}$ denoting $h_*^{'}$ denoting $h_4^{'}$ $(Q_4^*)$, and $[g_*^{'}]_4$

Denoting $[g^{'}]_4(P_4^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_4]_4 + [m_a]_4 = -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4 - [f^{'}{}_{4*}]_4 [p_4^*](a_4)M_4 - [f_{2*}^{'}]_4 p_4^*(a_4)W_4$$

$$m_4(t,0) = [h_*]_4 \int_0^\infty \beta_4(a_4)m_4(t,a_4)da_4 + \frac{[h_*^4]_4}{[h^*]_4}[p^*]_4(0)\, W_{,4} \qquad (30)$$

$$\frac{dW_4}{dt} = [g_*^4]_4 [Q_4^*]M_4$$

This is equation is important in the following cases: problem of using observed correlations to infer causal relations is relevant to (e) a wide variety of scientific disciplines. Yet given correlations between just two classical variables, (eb) it is impossible to determine whether they arose from (e) a causal influence of one on the other or a common cause influencing both, unless one can implement (eb) a randomized intervention. **Katja Ried, Megan Agnew,**

**Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. Resch** here consider the problem of causal inference for (e) quantum variables. Authors introduce causal tomography, which unifies and generalizes (e&eb) conventional quantum tomography schemes to provide (eb) a complete solution to the causal inference problem using (e) a quantum analogue of a randomized trial. Authors furthermore show that, in contrast to the classical case, observed quantum correlations alone can sometimes provide (eb) a solution. They implement a quantum-optical experiment that allows (eb) to control (e&eb) the causal relation between two optical modes, and two measurement schemes -- one with (e&eb) and one without (e) randomization -- that extract this relation from (e) the observed correlations. Results show that entanglement and (e&eb) coherence, known to be central to (e) quantum information processing, also provide (eb) a quantum advantage for (e) causal inference. Subjects: Quantum Physics (quant-ph); Learning (cs.LG); General Relativity and Quantum Cosmology (gr-qc); Machine Learning (stat.ML) Journal reference: Nat Phys 11, 414-420 (2015) DOI: 10.1038/nphys3266 Cite as: arXiv: 1406.5036 [quant-ph] (or arXiv: 1406.5036v1 [quant-ph] for this version) **Inferring causal structure: a quantum advantage Katja Ried, Megan Agnew, Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. ReschSTOP D. T. Son, A. O. Starinets** review recent progress in applying the AdS/CFT correspondence to (e&eb) finite-temperature field theory. In particular, they show how the hydrodynamic behavior of field theory is reflected in (e&eb) the low-momentum limit of (e) correlation functions computed through (e&eb) a real-time AdS/CFT prescription, which they formulate. Authors also show how the hydrodynamic modes in field theory correspond to (e&eb) the low-lying quasinormal modes of (e) the AdS black p-brane metric. They provide a proof of the universality of (e) the viscosity/entropy ratio within (eb) a class of theories with (e&eb) gravity duals and formulate (eb) a viscosity bound conjecture. Possible implications for real systems are mentioned. Subjects: High Energy Physics - Theory (hep-th) Journal reference: Ann.Rev.Nucl.Part.Sci.57:95-118,2007 DOI: 10.1146/annurev.nucl.57.090506.123120Report number: INT PUB 07-02Cite as: arXiv:0704.0240 [hep-th]

(or arXiv:0704.0240v2 [hep-th] for this version) An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society **Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP** Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv:

1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version **Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP Viscosity, Black Holes, and Quantum Field Theory D. T. Son, A. O. Starinets STOP Authors** study information retrieval from evaporating black holes, assuming that the internal dynamics of a black hole is unitary and rapidly mixing, and assuming that the retriever has unlimited control over the emitted Hawking radiation. If the evaporation of the black hole has already proceeded past the ``half-way" point, where (e) half of the initial entropy has been radiated away, then (eb) additional quantum information deposited in (eb) the black hole is revealed in (eb) the Hawking radiation very rapidly. Information deposited prior to the half-way point remains concealed until (e) the half-way point, and then emerges (eb) quickly. These conclusions hold because (e) typical local quantum circuits are efficient encoders for (e) quantum error-correcting codes that nearly achieve (eb) the capacity of the quantum erasure channel. Estimate of a black hole's information retention time, based on (e) speculative dynamical assumptions, is just barely compatible with (e&eb) the black hole complementarity hypothesis. **Black holes as mirrors: quantum information in random subsystems Patrick Hayden1 and John Preskill2 Published 26 September 2007 Journal of High Energy Physics, Volume 2007, JHEP09 (2007) STOP** By analyzing the key properties of black holes from the point of (e) view of quantum information, authors derive a model-independent picture of (e) black hole quantum computing. It has been noticed that this picture exhibits (eb) striking similarities with (e&eb) quantum critical condensates, allowing (eb) the use of a common language to describe (eb) quantum computing in both systems. Authors analyze such quantum computing by allowing coupling to (e&eb) external modes, under the condition that the external influence must be (=) soft-enough in order not to (e) offset the basic properties of the system. Authors derive model-independent bounds on some crucial time-scales, such as (=) the times of gate operation, decoherence, maximal entanglement and total scrambling. They show that for black hole type quantum computers all these time-scales are of (e) the order of the

black hole half-life time. Furthermore, authors construct explicitly a set of Hamiltonians that generates (eb) a universal set of quantum gates for (e) the black hole type computer. They find that the gates work at maximal energy efficiency. Furthermore, authors establish a fundamental bound on the complexity of quantum circuits encoded on (eb&e) these systems, and characterize the unitary operations that are (=) implementable. It becomes apparent that the computational power is very limited due to (e) the fact that the black hole life-time is of the same order of (e) the gate operation time. As a consequence, it is impossible to retrieve its information, within (eb) the life-time of a black hole, by externally coupling to (e&eb) the black hole qubits. However, authors show that, in principle, coupling to (e&eb) some of the internal degrees of freedom allows (eb) acquiring knowledge about the micro-state. Still, due to the trivial complexity of operations that can be performed, (eb) there is no time advantage over the collection of Hawking radiation and subsequent decoding. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Report number: LMU-ASC 03/16, MPP-2016-1 Cite as: arXiv:1605.01407 [hep-th] (or arXiv:1605.01407v1 [hep-th] for this version) **Universality of Black Hole Quantum Computing Gia Dvali, Cesar Gomez, Dieter Lust, Yasser Omar, and Benedikt Richter (Submitted on 4 May 2016) STOP** In this addendum to [arXiv: 1402.5674] two points are discussed. In the first additional evidence is provided (eb) for a dual connection between the geometric length of an Einstein-Rosen bridge and (e&eb) the computational complexity of the quantum state of the dual CFT's. The relation between growth of complexity and (e&eb) Page's “Extreme Cosmic Censorship” principle is also remarked on. The second point involves a gedanken experiment in which Alice measures (eb) a complete set of commuting observables at her end of (e) an Einstein-Rosen bridge is discussed. An apparent paradox is resolved by appealing to (e) the properties of GHZ tripartite entanglement. **Addendum to computational complexity and black hole horizons Leonard Susskind* Version of Record online: 18 JAN 2016 DOI: 10.1002/prop.201500093 STOP** Analyzing some well established facts, authors give a model-independent

parameterization of black hole quantum computing in terms of (e&*eb) a set of macro and micro quantities and their relations. These include the relations between the extraordinarily-small energy gap of (e) black hole qubits and important time-scales of (e) information-processing, such as, scrambling time and Page's time. Authors then show, confirming and extending previous results, that other systems of nature with (e&eb) identical quantum informatics features are (=) attractive Bose-Einstein systems at (eb) the critical point of quantum phase transition. Here authors establish a complete isomorphy between (e&eb) the quantum computational properties of these two systems. In particular, they show that the quantum hair of a critical condensate is (=) strikingly similar to (=) the quantum hair of a black hole. Irrespectively whether one takes the similarity between the two systems as a remarkable coincidence or as (=) a sign of a deeper underlying connection, the following is evident. Black holes are not unique in their way of (e) quantum information processing and they can manufacture black hole based quantum computers in labs by (e) taking advantage of quantum criticality. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Atomic Physics (physics. Atom-ph); Quantum Physics (quant-ph) Cite as: arXiv: 1601.01329 [hep-th] (or arXiv: 1601.01329v1 [hep-th] for this version) **Black Hole Based Quantum Computing in Labs and in the Sky Gia Dvali, Mischa Panchenko (Submitted on 6 Jan 2016) STOP Gia Dvali, Alexander Gußmann** show that the existence of black holes with (e&eb) classical skyrmion hair invalidates (e) standard proofs that global charges, such as (=) the baryon number, cannot be conserved by (e) a black hole. By carefully analyzing the standard arguments based on (e) a Gedankenexperiment in which a black hole is seemingly-unable to (e) return the baryon number that it swallowed, authors identify inconsistencies in this reasoning, which does not (e) take into the account neither the existence of skyrmion black holes nor the baryon/skyrmion correspondence. Authors then perform a refined Gedankenexperiment by incorporating (e) the new knowledge and show that (eb) no contradiction with conservation of baryon number takes place at (eb) any stage of black hole evolution. Analysis also indicates (eb) no conflict

between semi-classical black holes and (e&eb) the existence of baryonic gauge interaction arbitrarily-weaker than gravity. Next, authors study classical cross sections of a minimally-coupled massless probe scalar field scattered by (e) _ a skyrmion black hole. Authors investigate how the skyrmion hair manifests itself by (e) comparing this cross section with (e&eb) the analogous cross section caused by (e) a Schwarzschild black hole which has (e) the same ADM mass as (=) the skyrmion black hole. Here authors find an order-one difference in the positions of (e) the characteristic peaks in (eb) the cross sections. The peaks are shifted to smaller scattering angles when (e) the skyrmion hair is present. This comes from the fact that the skyrmion hair changes (e&eb) the near horizon geometry of the black hole when compared to (e&eb) a Schwarzschild black hole with same ADM mass. Authors keep the study of this second aspect general so that (e) the qualitative results obtained can also be applied to(e&eb) black holes with classical hair of different kind. Subjects: High Energy Physics - Theory (hep-th); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU-ASC 18/16 Cite as: arXiv:1605.00543 [hep-th] (or arXiv:1605.00543v2 [hep-th] for this version) A preliminary discussion is given of (e) the prospects that gravitational-wave observations of (e) binary inspiral of black holes could reveal or constrain (e&eb) quantum modifications to (e) black hole dynamics, such as are required to (e) preserve postulates of quantum mechanics. Different proposals for such modifications are characterized by (e) different scales, and the size of these scales relative to (e&eb) those probed by observation of (e) inspiral signals is important in determining (eb) the feasibility of finding experimental signatures. Certain scenarios with strong quantum modifications in (eb) a region extending well outside the horizon are expected to modify (e&eb) classical evolution, and distort (e&eb) the near-peak gravitational wave signal, suggesting (eb) a search for departures from (e) waveforms predicted by general relativity. The near agreement of the GW150914 signals with (e&eb) such waveforms is discussed, and indicates (eb) constraints on some such scenarios. Important strategies for more precise future tests are 1) to develop more precise

predictions from (e) scenarios proposing quantum modifications, and 2) searching for observed deviations from (e) numerical relativity predictions via (e&eb) analysis of gravity wave data, particularly focussing on (e&eb) the signal region corresponding to plunge and merger. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1602.03622 [gr-qc] (or arXiv: 1602.03622v2 [gr-qc] for this version) **Gravitational wave tests of quantum modifications to black hole structure -- with post-GW150914 update Steven B. Giddings STOP** Direct LHC bounds on (e&eb) colored SUSY particles now corner naturalness more than the measured value of the Higgs mass does. Bounds on the gluino are of (e) particular importance, since it radiatively "sucks" up (e) the stop and Higgs-up soft masses. As a result, even models that easily accommodate a 125 GeV Higgs are almost as tuned as (=) the simplest version of SUSY, the MSSM: at best at the percent level. In this paper, authors further examine how current LHC results constrain (e) naturalness in three classes of models that may relax LHC bounds on (e&eb) sparticles: split families (e&eb), baryonic RPV (e&eb), and Dirac gauginos. In models of split families and bRPV, the bounds on the gluino are (=) only slightly reduced, resulting in (eb) a few percent tuning. In particular, having a natural spectrum in bRPV models typically implies (eb) that tops, Ws, and Zs are easily produced (eb) in the cascade decays of (e) squarks and gluinos. The resulting leptons and missing energy push (eb) the gluino mass limit above 1 TeV. Even when the gluino has a Dirac mass and does not contribute to the stop mass at one loop, tuning reappears in calculable models because there is no symmetry imposing the supersoft limit. Authors conclude that, even if sparticles are found at LHC-14, naturalness will not emerge triumphant. Subjects: High Energy Physics - Phenomenology (hep-ph)DOI:10.1007/JHEP03(2014)022 Cite as:rXiv:1309.3568 [hep-ph] (or arXiv:1309.3568v2 [hep-ph] for this version) **The Last Vestiges of Naturalness Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang, Ken Van Tilburg, Giovanni Villadoro STOP** The lack of evidence for new physics beyond (e) the standard model at the LHC points to (e) a paucity of new particles near the weak scale. This suggests that the weak scale is tuned and that supersymmetry, if present at

all, is realized (eb) at higher energies. The measured Higgs mass constrains (e) the scalar sparticles to be below 10^5 TeV, while (e) gauge coupling unification favors (eb) Higgsinos below 100 TeV. Nevertheless, in many models gaugino masses are (=) suppressed and remain within reach of the LHC. Tuning the weak scale and (e&eb) the renormalization group evolution of (e) the scalar masses constrain (e) Split model building. Due to the small gaugino masses, either the squarks or the up-higgs often run tachyonic; in the latter case, successful electroweak breaking requires heavy higgsinos near the scalar sparticles. We discuss the consequences of tuning the weak scale and the phenomenology of several models of Split supersymmetry including anomaly mediation, U(1)_(B-L) mediation, and Split gauge mediation. Subjects: High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Experiment (hep-ex) DOI: 10.1007/JHEP02 (2013)126 Cite as: arXiv: 1210.0555 [hep-ph] (or arXiv: 1210.0555v2 [hep-ph] for this version) **Mini-Split Asimina Arvanitaki, Nathaniel Craig, Savas Dimopoulos, Giovanni Villadoro STOP** Light states associated with (e&eb) the hierarchy problem affect (e&eb) the Higgs LHC production and decays. Authors illustrate this within the MSSM and two simple extensions applying (e&eb) the latest bounds from (e) LHC Higgs searches. Large deviations in the Higgs properties are expected in (eb) a natural SUSY spectrum. The discovery of a non-Standard-Model Higgs may signal (eb) the presence of light stops accessible at (eb) the LHC. Conversely, the more the Higgs is Standard-Model-like, (eb) the more tuned the theory becomes. Taking the ratio of different Higgs decay channels at the LHC cancels (er&eb) the leading (eb) QCD uncertainties and potentially improves (eb) the accuracy in Higgs coupling measurements to (e) the percent level. This may lead to (eb) the possibility of doing precision Higgs physics at (eb) the LHC. Finally, authors entertain the possibility that the ATLAS excess around (e&eb) 125 GeV persists with (e&eb) a Higgs production cross-section that is enhanced compared to (e&eb) the SM. This increase can only be accommodated in extensions of (e) the MSSM and it may suggest (eb) that stops lie below 400 GeV, likely within (eb) reach of next year's LHC run. Subjects: High Energy Physics - Phenomenology (hep-ph) DOI: 10.1007/JHEP02 (2012)144 Cite

as: arXiv: 1112.4835 [hep-ph] (or arXiv: 1112.4835v2 [hep-ph] for this version) **A Non Standard Model Higgs at the LHC as a Sign of Naturalness Asimina Arvanitaki, Giovanni Villadoro STOP** It has recently been suggested that the presence of a plenitude of light axions, an Axiverse, is (=) evidence for the extra dimensions of string theory. Authors discuss the observational consequences of (e) these axions on astrophysical black holes through (e&eb) the Penrose superradiance process. When an axion Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e&eb) the black hole "nucleus" forming (eb) a gravitational atom in the sky. The occupation number of superradiant atomic levels fed by (e) the energy and angular momentum of the black hole grows exponentially. The black hole spins down and an axion Bose-Einstein condensate cloud forms (eb) around it. When the attractive axion self-interactions become (=) stronger than the gravitational binding energy, (eb) the axion cloud collapses, a phenomenon known in condensed matter physics as "Bosenova". The existence of axions is first diagnosed by (e) gaps in the mass vs (e&eb) spin plot of astrophysical black holes. For young black holes the allowed values of (e) spin are quantized, giving rise to (eb) "Regge trajectories" inside the gap region. The axion cloud can also be observed directly either through (e&eb) precision mapping of the near horizon geometry or through (e&eb) gravitational waves coming from (e) the Bosenova explosion, as well as axion transitions and (e&eb) annihilations in (eb) the gravitational atom. Estimates suggest (eb) that these signals are detectable in (eb) upcoming experiments, such as Advanced LIGO, AGIS, and LISA. Current black hole spin measurements imply (eb) an upper bound on (e&eb) the QCD axion decay constant of (e) 2 x 10^17 GeV, while (e) Advanced LIGO can detect (e&eb) signals from (e) a QCD axion cloud with (e&eb) a decay constant as low as (=) the GUT scale. Authors finally discuss the possibility of observing the gamma-rays associated with the Bosenova explosion and, perhaps, the radio waves from axion-to-photon conversion for the QCD axion. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Journal

reference:Phys.Rev.D83:044026, 2011 DOI: 10.1103/PhysRevD.83.044026 Cite as: arXiv: 1004.3558 [hep-th] (or arXiv: 1004.3558v2 [hep-th] for this version) **Exploring the String Axiverse with Precision Black Hole Physics Asimina Arvanitaki, Sergei Dubovsky (Submitted on 20 Apr 2010 (v1), last revised 16 Aug 2011 (this version, v2)) STOP Skyrmion Black Hole Hair: Conservation of Baryon Number by Black Holes and Observable Manifestations Gia Dvali, Alexander Gußmann STOP Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust** discuss BMS supertranslations both at (eb) null-infinity and on the horizon for the case of the Schwarzschild black hole. Authors show that both kinds of supertranslations lead to (eb) infinitely many gapless physical excitations. On this basis authors construct a quotient algebra using (e) suited superpositions of (e) both kinds of transformations which cannot be compensated by (e) an ordinary BMS-supertranslation and therefore are (=) intrinsically due to (e) the presence of an event horizon. Authors show that these quotient transformations are (=) physical and generate (eb) gapless excitations on the horizon that can account for (e) the gravitational hair as well as for (e) the black hole entropy. They identify the physics of these modes as associated with (e&eb) Bogolioubov-Goldstone modes due to quantum criticality. Classically the number of these gapless modes is (=) infinite. However, authors show that due to (e) quantum criticality the actual amount of information-carriers becomes (=) finite and consistent with Bekenstein entropy. Although they only consider the case of Schwarzschild geometry, the arguments are extendable to(e&eb) arbitrary space-times containing event horizons. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU--ASC 05/16, MPP--2016—3 Cite as: arXiv: 1601.03725 [hep-th] (or arXiv: 1601.03725v1 [hep-th] for this version) **Gravitational Black Hole Hair from Event Horizon Supertranslations Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust STOP Gia Dvali, Cesar Gomez, Nico Wintergerst** suggest that holography can be formulated in terms of (e&eb) the information capacity of (e) the St\"uckelberg degrees of freedom that maintain

gauge invariance of (e) the theory in the presence of (e) an information boundary. These St\"uckelbergs act as (e&eb) qubits that account for a certain fraction of (e) quantum information. Their information capacity is measured by (e) the ratio of the inverse St\"uckelberg energy gap to (e) the size of the system. Systems with (e&eb) the smallest gap are (=) maximally holographic. For massless gauge systems this information measure is (=) universally equal to (=) the inverse coupling evaluated at the systems' length scale. In this language it becomes very transparent why the St\"uckelberg information capacity of black holes saturates (e full) the Bekenstein bound and accounts for (e) the entire information of the system. The physical reason is that the strength of quantum interaction is bounded from below by the gravitational coupling, which scales as (=) area. Observing the striking similarity between the scalings of the energy gap of the boundary St\"uckelberg modes and (e&eb): (AC=DC=0) the Bogoliubov modes of critical many-body systems, authors establish a connection between holography and (e&eb) quantum criticality through (e&eb) the correspondence between these modes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1511.03525 [hep-th] (or arXiv: 1511.03525v1 [hep-th] for this version) **Stückelberg Formulation of Holography Gia Dvali, Cesar Gomez, Nico Wintergerst STOP B**lack hole entropy, denoted by N, in (semi)classical limit is (=) infinite. This scaling reveals (eb) a very important information about the qubit degrees of (e) freedom that carry black hole entropy. Namely, the multiplicity of qubits scales as N, whereas (e) their energy gap and their coupling as 1/N. Such a behavior is indeed exhibited by (e) Bogoliubov-Goldstone degrees of freedom of (e) a quantum-critical state of N soft gravitons (a condensate or a coherent state) describing (eb) the black hole quantum portrait. They can be viewed as (=) the Goldstone modes of (e) a broken symmetry acting on (e&eb) the graviton condensate. In this picture Minkowski space naturally emerges as (=) a coherent state of infinite-N gravitons of (e) infinite wavelength and it carries (e&eb) an infinite entropy. In this paper authors **Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015)** ask what is the geometric meaning (if any) of (e) the classical limit of this symmetry. Authors argue that the infinite-

N limit of Bogoliubov-Goldstone modes of critical graviton condensate is described by (e) recently-discussed classical BMS super-translations broken by the black hole geometry. However, the full black hole information can only be recovered for (e) finite N, since the recovery time becomes (=) infinite in classical limit in which N is (=) infinite. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) DOI: 10.1016/j.physletb.2015.11.073 Report number: MPP-2015-215, LMU-ASC 59/15 Cite as: arXiv: 1509.02114 [hep-th] (or arXiv: 1509.02114v1 [hep-th] for this version) **Classical Limit of Black Hole Quantum N-Portrait and BMS Symmetry Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015) STOP** Recent ideas about understanding physics of black hole information-processing in terms of (e&eb) quantum criticality allow us to (eb) implement black hole mechanisms of quantum computing within (eb) critical Bose-Einstein systems. The generic feature, uncovered both by analytic and numeric studies, is (=) the emergence at the critical point of (e) gapless weakly-interacting modes, which act as (=) qubits for information-storage at a very low energy cost. These modes can be effectively described in terms of (e&eb) either Bogoliubov or Goldstone degrees of freedom. The ground-state at the critical point is (=) maximally entangled and far from being classical. We confirm this near-critical behavior by a new analytic method. **Gia Dvali, Mischa Panchenko** compute growth of entanglement and show its consistency with (e&eb) black hole type behavior. On the other hand, in the over-critical regime the system develops (eb) a Lyapunov exponent and scrambles quantum information very fast. By, manipulating (e&eb) the system parameters externally, authors can put it in and out of various regimes and in this way control (e&eb) the sequence of information storage and processing. By using gapless Bogoliubov modes as (=) control qubits they design some simple logic gates. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Cite as: arXiv: 1507.08952 [hep-th] (or arXiv: 1507.08952v1 [hep-th] for this version) **Black Hole Type Quantum**

**Computing in Critical Bose-Einstein Systems Gia Dvali, Mischa Panchenko (Submitted on 31 Jul 2015) STOP Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst** establish a fundamental connection between quantum criticality of a many-body system, such as Bose-Einstein condensates, and (e&eb) its capacity of information-storage and processing. For deriving the effective theory of modes in the vicinity of the quantum critical point they develop a new method by mapping a Bose-Einstein condensate of N-particles onto (e&eb) a sigma model with a continuous global (pseudo)symmetry that mixes (e&eb) bosons of different momenta. The Bogolyubov modes of the condensate are mapped onto (e&eb) the Goldstone modes of the sigma model, which become (=) gapless at the critical point. These gapless Goldstone modes are (=) the quantum carriers of information and entropy. Analyzing their effective theory, authors observe the information-processing properties strikingly similar to (=) the ones predicted by the black hole portrait. The energy cost per qubit of information-storage vanishes in (eb) the large-N limit and the total information-storage capacity increases with (e&eb) N either exponentially or as a power law. The longevity of information-storage also increases with (e&eb) N, whereas (e) the scrambling time in (eb) the over-critical regime is controlled by (e&eb) the Lyapunov exponent and scales logarithmically with (e&eb) N. This connection reveals (eb) that the origin of black hole information storage lies in (eb) the quantum criticality of the graviton Bose-gas, and that much simpler systems that can be manufactured in table-top experiments can exhibit (eb) very similar information-processing dynamics. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 92, 125002 (2015) DOI: 10.1103/PhysRevD.92.125002 Cite as: arXiv:1507.02948 [hep-th] (or arXiv:1507.02948v1 [hep-th] for this version) **Nambu-Goldstone Effective Theory of Information at Quantum Criticality Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst STOP Gia Dvali, Cesar Gomez** show that a positive cosmological constant is (=) incompatible with the quantum-corpuscular resolution of de Sitter metric in form of (e) a coherent state. The reason is very general and is due to (e) the

quantum self-destruction of (e) the coherent state because of (e) the scattering of constituent graviton quanta. This process creates an irreversible quantum clock, which precludes (e) eternal de Sitter. It also eliminates (e) the possibility of Boltzmann brains and Poincare recurrences. This effect is expected to be (=) part of any microscopic theory that takes into account (e) the quantum corpuscular structure of the cosmological background. This observation puts the cosmological constant problem in a very different light, promoting it, from a naturalness problem, into (e&eb) a question of quantum consistency. We are learning that quantum gravity cannot tolerate (e) exceedingly-classical sources. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1412.8077 [hep-th]

(or arXiv:1412.8077v1 [hep-th] for this version) **Quantum Exclusion of Positive Cosmological Constant? Gia Dvali, Cesar Gomez (Submitted on 27 Dec 2014) STOP G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger** establish a connection between the ultra-Planckian scattering amplitudes in field and string theory and (e&eb) unitarization by black hole formation in these scattering processes. **Using as a guideline an explicit microscopic theory in which the black hole represents (eb) a bound-state of many soft gravitons at (eb) the quantum critical point, they were able to identify and compute (e&eb) a set of perturbative amplitudes relevant for (e) black hole formation. These are the tree-level N-graviton scattering S-matrix elements in (eb) a kinematical regime (called classicalization limit) where (e) the two incoming ultra-Planckian gravitons produce (eb) a large number N of soft gravitons. Authors compute these amplitudes by using (e) the Kawai-Lewellen-Tye relations, as well as scattering equations and (e&eb) string theory techniques. Authors discover that this limit reveals (eb) the key features of the microscopic corpuscular black hole N-portrait. In particular, the perturbative suppression factor of a N-graviton final state, derived from (e) the amplitude, matches the non-perturbative black hole entropy when (e) N reaches the quantum criticality value, whereas final states with different value of N are (=) either suppressed or excluded by non-perturbative**

**corpuscular physics. Thus they identify the microscopic reason behind the black hole dominance over other final states including (e) non-black hole classical object. In the parameterization of the classicalization limit the scattering equations can be solved exactly allowing (eb) us to obtain closed expressions for the high-energy limit of the open and closed superstring tree-level scattering amplitudes for a generic number N of external legs. Authors demonstrate matching and complementarity between the string theory and (e&eb) field theory** in different large-s and large-N regimes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1016/j.nuclphysb.2015.02.004 Report number: MPP-2014-320, LMU-ASC 52/14 Cite as: arXiv:1409.7405 [hep-th] (or arXiv:1409.7405v2 [hep-th] for this version) **Black Hole Formation and Classicalization in Ultra-Planckian 2 -> N Scattering G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger STOP** Gravitational backgrounds, such as black holes, AdS, de Sitter and inflationary universes, should be viewed as (=) composite of N soft constituent gravitons. It then follows that such systems are close to (e) quantum criticality of graviton Bose-gas to Bose-liquid transition. Generic properties of the ordinary metric description, including (e) geodesic motion or particle-creation in the background metric, emerge as (=) the large-N limit of quantum scattering of (e) constituent longitudinal gravitons. **Gia Dvali, Cesar Gomez** show that this picture correctly accounts for (e) physics of large and small black holes in AdS, as well as reproduces (eb) well-known inflationary predictions for cosmological parameters. However, it anticipates new effects not captured by (e) the standard semi-classical treatment. In particular, **Gia Dvali, Cesar Gomez** predict observable corrections that are sensitive to (e&eb) the inflationary history way beyond last 60 e-foldings. They derive an absolute upper bound on the number of e-foldings, beyond (e) which neither de Sitter nor inflationary Universe can be approximated by (e) a semi-classical metric. However, they could in principle persist in a new type of (e) it quantum eternity state. Authors discuss implications of this phenomenon for (e&eb) the cosmological constant problem. Subjects: High Energy Physics - Theory (hep-th); Cosmology and

Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Quantum Physics (quant-ph) DOI: 10.1088/1475-7516/2014/01/023 Cite as: arXiv: 1312.4795 [hep-th] (or arXiv:1312.4795v1 [hep-th] for this version) **Quantum Compositeness of Gravity: Black Holes, AdS and Inflation Gia Dvali, Cesar Gomez STOP** As quantum information science approaches the goal of constructing quantum computers, understanding loss of information through (e&eb) decoherence becomes (=) increasingly important. The information about a system that can be obtained from (e) its environment can facilitate quantum control and error correction. Moreover, observers gain (eb) most of their information indirectly, by (e) monitoring (primarily photon) environments of (e) the "objects of interest." Exactly how this information is inscribed in the environment is essential for (e) the emergence of "the classical" from (e) the quantum substrate. In this paper, **Robin Blume-Kohout, Wojciech H. Zurek** examine how many-qubit (or many-spin) environments can store (e) information about a single system. The information lost to (e) the environment can be stored redundantly, or it can be encoded in (eb) entangled modes of the environment. Authors go on to show that randomly chosen states of the environment almost always encode (e) the information so that (e) an observer must capture (e) a majority of the environment to deduce (eb) the system's state. Conversely, in the states produced by (e) a typical decoherence process, information about a particular observable of the system is stored redundantly. This selective proliferation of "the fittest information" (known as Quantum Darwinism) plays a key role in (eb) choosing the preferred, effectively classical observables of (e) macroscopic systems. The developing appreciation that the environment functions not just as (=) a garbage dump, but as a communication channel, is extending understanding of the environment's role in (e&eb) the quantum-classical transition beyond (e) the traditional paradigm of decoherence. Keywords Decoherence quantum Darwinism quantum information quantum-classical transition **Foundations of Physics November 2005, Volume 35, Issue 11, pp 1857-1876 A Simple Example of "Quantum Darwinism": Redundant Information Storage in Many-Spin Environments Robin Blume-Kohout, Wojciech H. Zurek STOP Harold Ollivier, David**

**Poulin, and Wojciech H. Zurek** study the emergence of objective properties in open quantum systems. In our analysis, the environment is promoted from a passive role of a reservoir selectively destroying quantum coherence to an active role of amplifier selectively **proliferating information about the system. We show that only preferred pointer states of the system can leave a redundant and therefore easily detectable imprint on the environment. Observers who—as is almost always the case—discover the state of the system indirectly (by probing a fraction of its environment) will find out only about the corresponding pointer observable. Many observers can act in this fashion independently and without perturbing the system. They will agree about its state. In this operational sense, preferred pointer states exist objectively. Received 24 July 2003 DOI:http://dx.doi.org/10.1103/PhysRevLett.93.220401 Objective Properties from Subjective Quantum States: Environment as a Witness Harold Ollivier, David Poulin, and Wojciech H. Zurek Phys. Rev. Lett. 93, 220401 – Published 22 November 2004 STOP** Quantum Darwinism recognizes that we—the observers—acquire our information about the "systems of interest" indirectly from their imprints on the environment. **Michael Zwolak, H. T. Quan, and Wojciech H. Zurek** show that information about a system can be acquired from a mixed-state, or hazy, environment, but the storage capacity of an environment fragment is suppressed by its initial entropy. In the case of good decoherence, the mutual information between the system and the fragment is given solely by the fragment's entropy increase. For fairly mixed environments, this means a reduction by a factor 1−h, where h is the haziness of the environment, i.e., the initial entropy of an environment qubit. Thus, even such hazy environments eventually reveal the state of the system, although now the intercepted environment fragment must be larger by ~ (1−h)−1 to gain the same information about the system. DOI:http://dx.doi.org/10.1103/PhysRevLett.103.110402 **Quantum Darwinism in a Mixed Environment Michael Zwolak, H. T. Quan, and Wojciech H. Zurek Phys. Rev. Lett. 103, 110402 – Published 8 September 2009** We study quantum Darwinism—the redundant recording of information about the preferred states of a decohering system by its environment—for an object illuminated by a

blackbody. In the cases of point-source and isotropic illumination, we calculate the quantum mutual information between the object and its photon environment. We demonstrate that this realistic model exhibits fast and extensive proliferation of information about the object into the environment and results in redundancies orders of magnitude larger than the exactly soluble models considered to date. DOI:http://dx.doi.org/10.1103/PhysRevLett.105.020404 **Quantum Darwinism in an Everyday Environment: Huge Redundancy in Scattered Photons C. Jess Riedel and Wojciech H. Zurek Phys. Rev. Lett. 105, 020404 – Published 9 July 2010 STOP** Measurements transfer information about a system to the apparatus and then, further on, to observers and (often inadvertently) to the environment. I show that even imperfect copying essential in such situations restricts possible unperturbed outcomes to an orthogonal subset of all possible states of the system, thus breaking the unitary symmetry of its Hilbert space implied by the quantum superposition principle. Preferred outcome states emerge as a result. They provide a framework for "wave-packet collapse," designating terminal points of quantum jumps and defining the measured observable by specifying its eigenstates. In quantum Darwinism, they are the progenitors of multiple copies spread throughout the environment—the fittest quantum states that not only survive decoherence, but subvert the environment into carrying information about them—into becoming a witness. DOI:http://dx.doi.org/10.1103/PhysRevA.76.052110 **Quantum origin of quantum jumps: Breaking of unitary symmetry induced by information transfer in the transition from quantum to classical Wojciech Hubert Zurek Phys. Rev. A 76, 052110 – Published 16 November 2007 STOP** The Born rule, a foundational axiom used to deduce probabilities of events from **wavefunctions, is indispensable in the everyday practice of quantum physics. It is also key in the quest to reconcile the ostensibly inconsistent laws of the quantum and classical realms, as it confers physical significance to reduced density matrices, the essential tools of decoherence theory. Following Bohr's Copenhagen interpretation, textbooks postulate the Born rule outright. However, recent attempts to derive it from (e) other quantum principles have been successful, holding promise for (e)**

**simplifying and clarifying the quantum foundational bedrock. A major family of derivations is based on envariance, a recently discovered symmetry of (e) entangled quantum states. Here, authors identify and experimentally test three premises central to (e) these envariance-based derivations, thus demonstrating, (eb) in the microworld, the symmetries from (e) which the Born rule is derived. Further, they demonstrate envariance in a purely local quantum system, showing (eb) its independence from relativistic causality.Subjects: Quantum Physics (quant-ph)Cite as:arXiv:1604.01471 [quant-ph] (or arXiv:1604.01471v1 [quant-ph] for this version) Quantum probabilities from quantum entanglement: Experimentally unpacking the Born rule Jérémie Harris, Frédéric Bouchard, Enrico Santamato, Wojciech H. Zurek, Robert W. Boyd, Ebrahim Karimi**

Integrating $[m_4]_4(t, a_4)$ with respect to a from 0 to ∞, we have

$$\frac{dM_4(t)}{dt} = \int_0^\infty [mt]_4(t, a_4)da_4$$

$$= -\int_0^\infty [m_a]_4\,(t, a_4)da_4 - \int_0^\infty \mu(a_4)m(t, a_4)da_4$$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f^{'}{}_{4*}]_4 \int_0^\infty p_4^*(a_4)da_4\right)M_4 - [f_{2*}^4]_4 \int_0^\infty p_4^*(a_4)da_4 W_4$$

$= m_4(t, 0) - \int_0^\infty \mu_4(a_4)\, m_4(t, a_4)da_4$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}^{i}]_4 P_4^*\right)M_4 - [f_{2*}^{i}]_4 P_4^* W_{\cdot 4}$$

From (30), we arrive at the equations

$$\frac{dW_4(t)}{dt} = \int_0^{\infty} ([h_*]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4$$

$$-\left([f_{4*}] + [f_{2*}]_4 + \left[f_{4*}^{'}\right]_4 P_4^*\right) M_4 + \left(\frac{\left[h_*^{'}\right]_4}{[h_*]^4}\, P_4^*(0) - f_{2*}^{'} P_4^*\right) W_4,$$

(34)

$$\frac{dW_4(t)}{dt}\left[g_{*}^{'}\right]_4 Q_4^*\, M_{.4}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{r*}]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4 = ([h_*]_4\beta_4(\theta_4) - \mu_4(\theta_4))M_{,4}$$

Where $0< \theta_4 = \theta_4(t) < \infty$.

And if $a_4^2\,[a_4^2]_4<4\,[a_2]_4<1\, a_2 e$,

$\mathrm{Re}\,\lambda_4\left[\lambda_4\right]_4 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_4 = -\frac{4}{2}a_4[a2\ ]_4 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_*\,[h_*]_4\beta_4(a_4) - \mu_4(a_4)\right|$. Then$|n_4(t)|$ ⊠⊿∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_4^2 <4\, a_2 e$ and $[a_4] > 2\Delta$,or $[a_4^2]$ ⊠[$a_2$]e and $(a_4 - \sqrt{(a_4^2)_4 - 4(a_2)_4 e}) > 2\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+4}\right)$ and $\beta_4(a_4) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4) + d\tau_4 da} = \frac{4}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_4 + [\mu_0]$ 0.5. Assume$[f_4](x_4)= K_4 x_4$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_4(x_4)= [0x_4]/(1 + x_4)$; the functional responses $[[f_2](x) = [0]\, x_4/(1 + x_4)$and $h_4(x_4)=(1+\omega_4 x_4)/(1 + x_4)$(See Holling [7] and Rashevsky [13])

Equation (44),now has the form of

$$\frac{m_4 + [\mu_0] + [K][P_0])}{[m_4]_4 + (\mu_0)_4 + [K][P_0])\frac{3}{4}\ [K]_4 [P_0]\frac{3}{4}} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(4 + P_0) < b_4$ i e $[P_0] < b/(\delta_4 - b_4)$then the equilibrium$(P_0)$ $[P_0]_4(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_4$

$[f_{4*}] = K_4\ P^*\ [f_{2*}] = \frac{\sigma Q_4^*}{4+Q_4^*}$ $[h_*] = \frac{4+\omega_4 Q_4^*}{4+Q_4^*}$

$[f_{4*}] = K_4\ [f^4_{2*}] = \frac{\sigma_4}{(4+Q_4^*)^2}[\ g^4_*] = \frac{\sigma_4}{(4+P_4^*)^2}$ $[h^4_*] = \frac{\omega_4 - 4}{(4+Q_4^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0$, K, b, $\delta$, and $\omega$ need to satisfy

$$\frac{\omega_4}{(m + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma)^2} + \frac{2\omega_4}{([m_4]_4 + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma_4)^3} < 1$$

And

$$\frac{1}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^2}+\frac{2\omega_4}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4}+(m_4+[\mu_0]_4+\frac{K_4b_4}{\delta_4-b_4}+\sigma_4)\,Q_4^*)^3$$

$$=m_4+[\mu_0]_4+3+\frac{K_4b_4}{\delta_4-b_4}(4+Q^*)(1+\omega_4Q_4^*$$

Now

$$\left[a_4^2\right]-1\left[a_2\right]e=\left(\frac{2K_4b_4}{\delta_4-b_4}\right)^2+\frac{1K_4b_4\sigma_4\,Q_4^*}{(\delta_4-b_4)(1+Q_4^*)}+\frac{\sigma_4^2Q_4^{*2}}{(1+Q_4^*)^2}$$

$$\frac{1\,b_4\sigma_4\,(\delta_4-b)Q^*}{\delta_4(1+Q_4^*)^2}-\frac{1\,(\delta_4-b_4)^2(1-\omega_4)Q_4^*}{\delta_4(1+\omega_4Q_4^*)(1+Q_4^*)}P^*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_4\,ae_4^{mu}\;[\mu_0]_4-\frac{a_4}{a_4+4}$$

Hence if $[a_4^2>4(a_2)_4e$ and

$$[a_4] = \frac{2K_4 b_4}{\delta_4 - b_4} + \frac{\sigma_4 \, Q_4^*}{1 + Q_4^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_4]_4 = [f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}] P_4^*$$

$$[a_2]_4 = \left[f_{2*}^4\right] P_4^* - \frac{h_*^4}{h_*} P^*(0)$$

$$e = [g_*^4]_4 \, Q_4^*$$

$$\eta_4(t) = \left[h_*\right]_4 \beta_4 \left(\theta_4(t)\right) - \mu_4(\theta(t)),$$

The system(34)can be written as the following system:

$$[M(t)]\left[a_4\right] M_4 - [a_2]_4 W_4 - [n(t)] M_4$$

$$w_4(t) = eM_4$$

Write

$$A = \begin{pmatrix} -[a_4]_4 - [a_2]_4 & \\ e & 0 \end{pmatrix}$$

$$B = (t) = \begin{pmatrix} [\eta(t)]_4 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_4 \\ W_4 \end{pmatrix}$$

The solution of (32)is

$$X_4(t) = e^{A,t}[X]_4(0) + \int_0 e^{At}$$

**Some of the other conceptual and thematic and discursive factorial representationalities to which the model is applicable are:every one-dimensional quantum mechanical Hamiltonian H1 can have (e) a partner H2 such that H1 and H2 taken together may be viewed as (=) the components of a supersymmetric Hamiltonian. The term 'supersymmetric Hamiltonian' is taken to mean (eb) a Hamiltonian defined in terms of (e&eb)charges that obey (eb)the same algebra as that of (e) the generators of supersymmetry in field theory. The consequences of this symmetry for (er) the spectra of H1 and H2 are explored. It is shown how the supersymmetric pairing may be utilised to eliminate (e) the ground state of H1, or add a state below the ground state of H1 or maintain the spectrum of H1. It is also explicitly demonstrated that the supersymmetric pairing may be used to generate (eb) a class of anharmonic potentials with (e&eb) exactly specified spectra. The complete spectrum of an anharmonic potential so generated consists of (e) all the eigenstates of the simple harmonic oscillator and, in addition, a ground state at a specified energy E which lies arbitrarily below (eb) the E=1/2 ground state of the harmonic oscillator. Supersymmetric quantum mechanics of one-dimensional systems C V Sukumar Journal of Physics A: Mathematical and General, Volume 18, Number 15 STOP Authors examine the properties and consequences of (e&eb) pseudo-supersymmetry for (e) quantum systems admitting (e) a pseudo-Hermitian Hamiltonian. Authors explore the Witten index of (e) pseudo-**

**supersymmetry and show that every pair of diagonalizable (not necessarily Hermitian) Hamiltonians with (e&eb) discrete spectra and (e&eb) real or complex-conjugate pairs of (e) eigenvalues are (=) isospectral and have (e) identical degeneracy structure except (e) perhaps for the zero eigenvalue if and only if (e) they are pseudo-supersymmetric partners. This implies (eb) that pseudo-supersymmetry is (=) the basic framework for (e) generating non-Hermitian PT-symmetric and (e&eb) non-PT-symmetric Hamiltonians with (e&eb) a real spectrum via (e&eb) a Darboux transformation, and shows (eb) that every diagonalizable Hamiltonian H with(e&eb) a discrete spectrum and (e&eb) real or complex-conjugate pairs of (e) eigenvalues may be factored as (=) H=L♯L where L is (=) a linear operator with pseudo-adjoint L♯. In particular, this factorization applies to (e&eb) PT-symmetric and Hermitian Hamiltonians. The non-degenerate two-level systems provide (eb) a class of Hamiltonians that are pseudo-Hermitian. Authors demonstrate the implications of general results for (e&eb) this class in some detail. Copyright © 2002 Elsevier Science B.V. All rights reserved. Nuclear Physics B Volume 640, Issue 3, 30 September 2002, Pages 419–434Pseudo-supersymmetric quantum mechanics and isospectral pseudo-Hermitian Hamiltonians Ali Mostafazadeh STOP Michele Caselle, Alessandro Nada, Marco Panero present a high-precision lattice calculation of (e) the equation of state in the confining phase of (e) SU(2) Yang-Mills theory. Authors show that the results are described very well by (e) a gas of massive, non-interacting glueballs, provided (e) one assumes an exponentially growing Hagedorn spectrum. The latter can be derived within (eb) an effective bosonic closed-string model, leading to (eb) a parameter-free theoretical prediction, which is in perfect agreement with (=) proposed lattice results. Furthermore, when applied to (e&eb) SU (3) Yang-Mills theory, this effective model accurately describes (eb) the lattice results reported by Bors\'anyi et al. in JHEP 07 (2012) 056. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 07 (2015) 143 DOI: 10.1007/JHEP07 (2015)143 Cite as: arXiv:**

**1505.01106 [hep-lat] (or arXiv: 1505.01106v2 [hep-lat] for this version) Hagedorn spectrum and thermodynamics of SU (2) and SU (3) Yang-Mills theories Michele Caselle, Alessandro Nada, Marco Panero STOP B. de Wit J. Hoppe, H. Nicolai study the quantum-mechanical properties of (e) a supermembrane and examine the nature of (e) its ground state. A supersymmetric gauge theory of (e) area-preserving transformations provides (eb) a convenient framework for this study. The supermembrane can be viewed as (=) a limiting case of a class of models in supersymmetric quantum mechanics. Its mass does not depend on (=) the zero modes and vanishes only if (e) the wave function is (=) a singlet under (e&eb) supersymmetry transformations of (e) the nonzero modes. Authors exhibit (eb) the complexity of (e) the supermembrane ground state and (e&eb) examine various truncations of these models. None of these truncations has massless states. Copyright © 1988 Published by Elsevier B.V. Nuclear Physics B Volume 305, Issue 4, 28 December 1988: Pages 545-581: On the quantum mechanics of supermembranes Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. B. de Wit J. Hoppe, H. Nicolai STOP In this note we explain how world-volume geometries of (e) D-branes can be reconstructed within (eb) the microscopic framework where (e) D-branes are described through (e&eb) boundary conformal field theory. Authors extract the (non-commutative) world-volume algebras from (e) the operator product expansions of (e) open string vertex operators. For branes in a flat background with (e&eb) constant non-vanishing B-field, the operator products are computed perturbatively to (e) all orders in the field strength. Resulting series coincides with (=) Kontsevich's presentation of the Moyal product. After extending these considerations to (e) fermionic fields authors conclude with some remarks on the generalization of approach to (e&eb) curved backgrounds. D-branes and deformation quantization Volker Schomerus1,2 Published 25 July 1999 Journal of High Energy Physics, Volume 1999, JHEP06(1999)STOP Authors employ the path integral approach developed in Gangopadhyay S. and Scholtz F. E., Phys. Rev.**

**Lett., 102 (2009) 241602 to discuss the (generalized) harmonic oscillator in (eb) a noncommutative plane. The action for (e) this system is derived in (e&eb) the coherent state basis with (e&eb) additional degrees of freedom. From this the action in the coherent state basis without (e) any additional degrees of freedom is obtained. This gives (eb) the ground-state spectrum of the system. Authors then employ the exact renormalization group approach to show (eb) that equivalence can be constructed between this (noncommutative) system and (e&eb) a commutative system Path integral action and exact renormalization group dualities for quantum systems in noncommutative plane Sunandan Gangopadhyay Published 25 June 2015 • Copyright © EPLA, 2015 STOP Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini present a lattice study of (e)the equation of state in Yang-Mills theory based on (e) the exceptional G(2) gauge group. As is well-known, at zero temperature this theory shares (e&eb) many qualitative features with real-world QCD, including (e) the absence of colored states in the spectrum and dynamical string breaking at (e) large distances. In agreement with previous works, authors show that at finite temperature this theory features (eb) a first-order deconfining phase transition, whose nature can be studied by (e) a semi-classical computation. Authors also show that the equilibrium thermodynamic observables in the deconfined phase bear striking quantitative similarities with (=) those found in SU (N) gauge theories: in particular, these quantities exhibit (eb) nearly perfect proportionality to the number of gluon degrees of freedom, and (e&eb) the trace anomaly reveals (eb) a characteristic quadratic dependence on (e&eb) the temperature, also observed in SU (N) Yang-Mills theories (both in four and in three spacetime dimensions). Authors compare our lattice data with (e&eb) analytical predictions from (e) effective models, and discuss their implications for (e) the deconfinement mechanism and (e&eb) high-temperature properties of (e) strongly interacting, non-supersymmetric gauge theories. Results give (eb) strong evidence for (e) the conjecture that the thermal deconfining transition is governed by (e&eb) a universal mechanism, common to (e) all simple gauge groups.**

**Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 03 (2015) 057 DOI: 10.1007/JHEP03 (2015)057 Report number: DESY-14-146; IFT-UAM/CSIC-14-076 Cite as: arXiv: 1409.8305 [hep-lat] (or arXiv: 1409.8305v2 [hep-lat] for this version Exceptional thermodynamics: The equation of state of G(2) gauge theory Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini STOP In U (1) lattice gauge theory in three spacetime dimensions, the problem of confinement can be studied analytically in (e&eb) a semi-classical approach, in terms of (e&eb) a gas of monopoles with (e&eb) Coulomb-like interactions. In addition, this theory can be mapped to (e&eb) a spin model via (e&eb) an exact duality transformation, which allows (eb) one to perform high-precision numerical studies of (e) the confining potential. Taking advantage of these properties, authors carried out an accurate investigation of (e) the effective string describing (eb) the low-energy properties of flux tubes in (eb) this confining gauge theory. Authors found striking deviations from (e) the expected Nambu-Goto-like behavior, and, for the first time, evidence for (e) contributions that can be described by a term proportional to the extrinsic curvature of the effective string worldsheet. Such term is allowed by (e) Lorentz invariance and its presence in the infrared regime of the U (1) model was indeed predicted by (e) Polyakov several years ago. Results show that this term scales as expected according to (e) Polyakov's solution, and becomes (=) the dominant contribution to (e) the effective string action in the continuum limit. Authors also demonstrate analytically that the corrections to (e) the confining potential induced by (e) the extrinsic curvature term can be related to (e&eb) the partition function of (e) the massive perturbation of (e) a c=1 bosonic conformal field theory. The implications of results for SU (N) Yang-Mills theories in (e&eb) three and in four spacetime dimensions are discussed. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th) Journal reference: JHEP 1501 (2015) 105 DOI: 10.1007/JHEP01 (2015)105 Report number: IFT-UAM/CSIC-14-052 Cite as: arXiv:**

**1406.5127 [hep-lat] (or arXiv: 1406.5127v2 [hep-lat] for this version) A different kind of string Michele Caselle, Marco Panero, Roberto Pellegrini, Davide Vadacchino STOP In the past ten years, the ideas of supersymmetry have been profitably applied to (e&eb) many nonrelativistic quantum mechanical problems. In particular, there is now a much deeper understanding of why certain potentials are analytically solvable and (e&eb) an array of powerful new approximation methods for (e) handling potentials which are not (e) exactly solvable. In this report, authors review the theoretical formulation of (e) supersymmetric quantum mechanics and discuss many applications. Exactly solvable potentials can be understood in terms of (e&eb) a few basic ideas which include (e) supersymmetric partner potentials (e&eb) shape invariance and (e&eb) operator transformations. Familiar solvable potentials all have (e) the property of shape invariance. Authors describe new exactly solvable shape invariant potentials which include (e) the recently discovered self-similar potentials as (=) a special case. The connection between inverse scattering,(e&eb) isospectral potentials and (e&eb) supersymmetric quantum mechanics is discussed and multisoliton solutions of (e) the K dV equation are constructed. Approximation methods are also discussed within (eb) the framework of supersymmetric quantum mechanics and in particular it is shown that a supersymmetry inspired WKB approximation is (=) exact for a class of shape invariant potentials. Supersymmetry ideas give (eb) particularly nice results for the tunneling rate in (eb) a double well potential and for improving (eb+) large N expansions. Authors also discuss the problem of (e) a charged Dirac particle in (eb) an external magnetic field and other potentials in terms of (e&eb) supersymmetric quantum mechanics. Finally, authors discuss structures more general than supersymmetric quantum mechanics such as(=) parasupersymmetric quantum mechanics in which there is symmetry between a boson and (e&eb) a para-fermion of order p. Copyright © 1995 Published by Elsevier B.V. Physics Reports Volume 251, Issues 5–6, January 1995, Pages 267–385Supersymmetry and quantum mechanics Fred Coopera, Avinash Khareb, Uday Sukhatmec STOP The free energy of (e) the maximally**

**supersymmetric SU (N) gauge theory at temperature T is expected to scale, in the large-N limit, as (=) N2T4 times a function of the 't Hooft coupling, f (gYM2N). In the strong coupling limit the free energy has been deduced from (e) the near-extremal 3-brane geometry, and its normalization has turned out to be Full-size image (<1 K) times that found in (eb) the weak coupling limit. In this paper authors calculate the leading correction to (e) this result in inverse powers of (e) the coupling, which originates from (e) the R4 terms in the tree-level effective action of (e) type IIB string theory. The correction to Full-size image (<1 K) is positive and of (e) order Full-size image (<1 K). Thus, f (gYM2N) increases as (e&eb) the 't Hooft coupling is decreased, in accordance with (e&eb) the expectation that is should be approaching 1 in (eb) the weak coupling limit. Authors also discuss similar corrections for (e) other conformal theories describing (eb) coincident branes. In particular, authors suggest (eb) that the coupling independence of (e) the near extremal entropy for (e) D1-branes bound to (e) D5-branes is related to (e&eb) the vanishing of (e) the Weyl tensor of AdS3 × S3. Copyright © 1998 Published by Elsevier B.V. Nuclear Physics B Volume 534, Issues 1–2, 23 November 1998, Pages 202–222Coupling constant dependence in the thermodynamics of N = 4 supersymmetric Yang-Mills theory Steven S. Gubsera, Igor R. Klebanovb, 1, Arkady A. Tseytlinc, 2 STOP A new approach to (e) N=2 supersymmetry based on (e) the concept of harmonic superspace is proposed and is used to give (eb) an unconstrained superfield geometric description of N=2 super Yang-Mills and supergravity theories as (=) well as of matter N=2 hypermultiplets. The harmonic N=2 superspace has (e) an independent coordinates, in addition to the usual ones, the isospinor harmonics ui+or- on (e&eb) the sphere SU (2)/U (1). The role of ui+or- is to relate the SU (2) group realised on the component fields to (e&eb) a U (1) group acting on (e&eb) the relevant superfields. Their introduction makes it possible to SU (2)-covariantise (eb (t), e&eb) the notion of Grassmann analyticity. Crucial for the construction is the existence of (e) an analytic subspace of the general harmonic N=2 superspace. The hypermultiplet superfields and (e&eb) the true prepotentials (pre-**

**prepotentials) of N=2 super Yang-Mills and supergravity are (=) unconstrained superfunctions over (e&eb) this analytic subspace. The pre-prepotentials have (e) a clear geometric interpretation as gauge connections with (e&eb) respect to the internal SU (2)/U (1) directions. A radically new feature arises: the number of gauge and (e&eb) auxiliary degrees of freedom becomes (=) infinite while (e) the number of physical degrees of freedom remains (=) finite. Other new results are the massive N=2 Yang-Mills theory and (e&eb) various off-shell self-interactions of hypermultiplets. The propagators for matter and Yang-Mills superfields are given. Unconstrained N=2 matter, Yang-Mills and supergravity theories in harmonic superspace A Galperin, E Ivanov, S Kalitzin, V Ogievetsky and E Sokatchev STOP By assuming the existence of a novel multipronged string state for (e) D-particles interacting with (e&eb) D-brane intersections in type IIA string theory, authors are able to derive a quantum mechanical description of (e) supersymmetric Reissner-Nordstrom black holes. A supersymmetric index calculation provides (eb) evidence for (e) this conjecture. The quantum mechanical system becomes (=) two decoupled conformal quantum mechanical systems in (eb) the low energy limit. The conformal quantum mechanics has expected properties of (e) a dual description of string theory on (e&eb) AdS2×S2. Subjects: High Energy Physics - Theory (hep-th) Journal reference: JHEP 0001:006, 2000 DOI: 10.1088/1126-6708/2000/01/006 Report number: LBNL-43734 Cite as: arXiv:hep-th/9907101 (or arXiv:hep-th/9907101v5 for this version) Supersymmetric Quantum Mechanical Description of Four Dimensional Black Holes Julie D. BlumSTOP Using recently proposed soliton equations authors derive a basic identity for the scaling violation of (e) N=2 supersymmetric gauge theories Σiai∂F/∂ai−2F=8πib1u. Here F is (=) the prepotential, ai's are (=) the expectation values of (e) the scalar fields in (eb) the vector multiplet, u=1/2 Tr <ϕ2> and b1 is the coefficient of (e) the one-loop β-function. This equation holds in (eb) the Coulomb branch of all N=2 supersymmetric gauge theories coupled with (e&eb) massless matter. TOHRU EGUCHI and SUNG-KIL YANG, Mod Phys. Lett. A 11, 131 (1996).**

**DOI: 10.1142/S0217732396000151 PREPOTENTIALS OF N=2 SUPERSYMMETRIC GAUGE THEORIES AND SOLITON EQUATIONS STOP The Coulomb branch of N = 2 supersymmetric gauge theories in four dimensions is described in general by (e) an integrable Hamiltonian system in (eb) the holomorphic sense. A natural construction of such systems comes from (e) two-dimensional gauge theory and (e&eb) spectral curves. Starting from this point of view, authors propose an integrable system relevant to (e&eb) the N = 2 SU (n) gauge theory with (e&eb) a hypermultiplet in the adjoint representation, and offer (eb) much evidence that it is correct. The model has (e) a Full-size image (<1 K) S-duality group (with the central element −1 of Full-size image (<1 K) acting as (=) charge conjugation); Full-size image (<1 K) permutes the Higgs, confining, and oblique confining phases in the expected fashion. We also study more exotic phases. Copyright © 1996 Published by Elsevier B.V. Nuclear Physics B Volume 460, Issue 2, 5 February 1996, Pages 299–334Supersymmetric Yang-Mills theory and integrable systems Ron Donagia, 1, Edward Wittenb, 2 STOP A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 present a first step towards generalizing (eb) the work of Seiberg and Witten on N = 2 supersymmetric Yang-Mills theory to (e&eb) arbitrary gauge groups. Specifically, authors propose a particular sequence of (e) hyperelliptic genus n − 1 Riemann surfaces to underly (e&eb) the quantum moduli space of SU (n)N = 2 supersymmetric gauge theory. These curves have (e) an obvious generalization to (e) arbitrary simply laced gauge groups, which involves (e&eb) the A-D-E type simple singularities. To support proposal, authors argue that the monodromy in (e&eb) the semiclassical regime is (=) correctly reproduced. We also give some remarks on a possible relation to string theory. Copyright © 1995 Published by Elsevier B.V. Physics Letters B Volume 344, Issues 1–4, 26 January 1995, Pages 169–175Simple singularities and N = 2 supersymmetric Yang-Mills theory A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 STOP Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 review and extend earlier work that uses (e) the AdS/CFT correspondence to relate the**

**black-hole–black-string transition of gravitational theories on a circle to (e&eb) a phase transition in maximally supersymmetric (1 + 1)-dimensional SU(N) gauge theories at (eb) large N, again compactified on (e&eb) a circle. Authors perform gravity calculations to determine (eb) a likely phase diagram for (e) the strongly coupled gauge theory. Authors then directly study the phase structure of (e) the same gauge theory, now at (eb) weak't Hooft coupling. In the interesting temperature regime for (e) the phase transition, the (1 + 1)-dimensional theory reduces to (e&eb) a (0 + 1)-dimensional bosonic theory, which they solve using (e) Monte Carlo methods. Authors find strong evidence that the weakly coupled gauge theory also exhibits (eb) a black hole–black string-like phase transition in (eb) the large N limit. Authors demonstrate that a simple Landau–Ginzburg-like model describes (eb) the behaviour near the phase transition remarkably well. The weak coupling transition appears to be (=) close to the cusp between a first-order and a second-order transition. Black-hole–black-string phase transitions in thermal (1 + 1)-dimensional supersymmetric Yang–Mills theory on a circle Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 Published 29 October 2004 • 2004 IOP Publishing Ltd Classical and Quantum Gravity, Volume 21, Number 22STOP Monitoring the orbits of stars around (e&eb) Sgr A* offers(eb)the possibility of detecting (e&eb) the precession of (e) their orbital planes due to (e) frame dragging, of (e) measuring the spin and quadrupole moment of (e) the black hole, and of (e)testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and (e&eb) tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are (=), in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for which an optical interferometer, such as GRAVITY, will detect orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to (e&eb) the timescale of precession due to (e) the quadrupole moment of (e) the black hole. As a result, this non-conservative effect is (=) a**

**potential source of systematic uncertainty in (eb) testing the no-hair theorem with (e&eb) stellar orbits. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI:10.1088/0004-637X/777/1/57 Cite as:arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this version) Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP Authors perform a phase space analysis of (e) a generalized modified gravity theory with (e&eb) non-minimal coupling between geometry and (e&eb) matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that (eb) in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, Authors prove that there is (=) an attractor solution to describe (eb) the late time accelerating universe when (e) the modified gravity is chosen in (eb) a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and supernova type Ia data sets, Authors find that (eb) the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as:arXiv: 1512.09281 [gr-qc] (or arXiv: 1512.09281v1 [gr-qc] for this version) Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang study the P−V criticality and phase transition in (eb) the extended phase space of (e) charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of the system and (e&eb) its conjugate quantity is (=) the thermodynamic volume of (e) the black hole. The black**

**holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb)no P−V criticality and phase transition appear, while for the black holes with a spherical horizon, even when (e) the charge of the black hole is absent, the P−V criticality and the small black hole/large black hole phase transition (eb) will appear, but it happens only in (eb) d=5 dimensions; when (e) the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to (e) have the P−V criticality and the small black hole/large black hole phase transition, there exists (eb) an upper bound for (e) the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the Gauss-Bonnet coefficient and Q is the charge of the black hole. Authors calculate the critical exponents at the critical point and find that for all cases, they are (=) the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09 (2013)005 Cite as: arXiv: 1306.6233 [gr-qc] (or arXiv: 1306.6233v4 [gr-qc] for this version) P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP paper studies the whole process of (e) gravitational collapse and (e&eb) accretion of (e) a massless scalar field in (eb) asymptotically flat spacetime. Two kinds of initial configurations are considered. One is the initial data without (e) black hole; the other contains (e) a black hole. Under suitable initial conditions, Authors find that multi-horizon will appear which mean (eb) that the initial black hole formed by (e) gravitational collapse or existing at the beginning will instantly expand and suddenly grow (e&eb, eb), rather than grows gradually in the accretion process. A new type of critical behavior is found around (e&eb) the instant expansion. The numerical computation shows (eb) that the critical exponents are (=) universal. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Computational Physics (physics. Comp-ph) Cite as: arXiv: 1604.03363 [gr-qc] (or arXiv: 1604.03363v1 [gr-qc] for this version)**

**Multi-horizon and Critical Behavior in Gravitational Collapse of Massless Scalar Zhoujian Cao, Rong-Gen CAI, Run-Qiu Yang STOP In this paper, Authors present the results on the gravitational collapse of (e) charged massless scalar field in asymptotically flat spacetime with(e&eb) a perfectly reflecting wall. Differing from previous works, Authors study the system in (eb) the double null coordinates, by which we could simulate (e&eb) the system until the black hole forms with (e&eb) higher precision but less performance time. Authors investigate the influence of charge on (e&eb) the black hole formation and the scaling behavior near the critical collapses. The gapless and gapped critical behaviors for (e) black hole mass and charge are studied numerically. Authors find that they satisfy (eb) the scaling laws for critical gravitational collapse but the gapped critical behavior is different from (e) its AdS counterpart. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1602.00112 [gr-qc] (or arXiv: 1602.00112v1 [gr-qc] for this version) Multiple critical gravitational collapse of charged scalar with reflecting wall Rong-Gen CAI, Run-Qiu Yang STOP In this paper Authors investigate the (2+1)-dimensional topological non-Fermi liquid in (eb) strongly correlated electron system, which has (e) a holographic dual description by (e) Einstein gravity in (3+1)-dimensional anti-de Sitter (AdS) space-time. In a dyonic Reissner-Nordstrom black hole background, Authors consider a Dirac fermion coupled to (e&eb) the background U (1) gauge theory and an (e&eb) intrinsic chiral gauge field bM induced by (e) chiral anomaly. UV retarded Green's function of (e) the charged fermion in (eb) the UV boundary from (e) AdS4 gravity is calculated, by imposing (e&eb) in-falling wave condition at the horizon. Authors also obtain IR correlation function of (e) the charged fermion at (eb) the IR boundary arising from (e) the near horizon geometry of (e) the topological black hole with (e&eb) index k=0,±1. By using (e) the UV retarded Green's function and IR correlation function Authors analyze the low frequency behavior of (e) the topological non-Fermi liquid at zero and (e&eb) finite temperatures, especially the relevant non-Fermi liquid behavior near (eb) the quantum critical region. In addition, Authors**

**find that when k=±1, the effective mass of (e) bulk Dirac fermion at IR fixed point is (=) topologically quantized, consequently the Fermi momentum presents (eb) a discrete sparse distributed pattern, in contrary to (e&eb) the Ricci flat case with k=0. As demonstrated examples, Authors calculate the spectral functions of (e) topological non-Fermi liquid due to (e) anomalous Hall effect (AHE) (|b→|≠0,b 0=0) and chiral magnetic effect (CME) (|b→|=0,b0≠0), which can be tested in topological materials such as (=) topological insulator and (e&eb) topological metal with (e&eb) strong couplings. Subjects: High Energy Physics - Theory (hep-th); Strongly Correlated Electrons (cond-mat.str-el); General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv: 1601.03865 [hep-th] (or arXiv: 1601.03865v1 [hep-th] for this version) Topological Non-Fermi Liquid Rong-Gen Cai, Yong-Hui Qi, Yue-Liang Wu, Yun-Long Zhang STOP The unusual magnetic materials are significant in both science and technology. However, because of (e) the strongly correlated effects, it is difficult to understand their novel properties from (e) theoretical aspects. Holographic duality offers (eb) a new approach to understanding such systems from (e) gravity side. This paper will give a brief review of Authors recent works on the applications of holographic duality in understanding (e&eb) unusual magnetic materials. Some quantitative compare between holographic results and (e&eb) experimental data will be shown and some predictions from (e) holographic duality models will be discussed. Comments: For Second LeCosPA Symposium at Taiwan National University, Taipai, Dec. 14-18, 2015 Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1601.02936 [hep-th] (or arXiv:1601.02936v1 [hep-th] for this version) Understanding strongly coupling magnetism from holographic duality Rong-Gen Cai, Run-Qiu Yang STOP It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a**

**Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3**

**[astro-ph.CO] for this version) Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang**

# Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

## Module Number Six

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-465) 1990. We demonstrate its applicability cosmology/ quantum computation economics etal. For the first time, physicists couple distant nuclear spins using a single electron July 11, 2016. Some sentences have been deleted due to spatial restrictions. Please excuse me. For the first time, researchers at the University of Basel in Switzerland have coupled the nuclear spins of distant atoms using (e) just a single electron. Three research groups took part in this complex experiment, the results of which have now been published in the journal Nature Nanotechnology. In most materials, the nuclear spins of (e) neighboring atoms have (e) only a very weak effect on one another, as (=) the tiny nuclei are located deep within (e&eb) the atoms. This is not the case in metals, however, where (e) some of the electrons are able to move freely. The electron spins are able to couple nuclear spins at (eb) relatively large distances from one another. Named after four physicists, this RKKY interaction was discovered in the 1950s. Nuclear spins linked by (e&eb) an individual electron spin Now, an experiment by researchers from the Department of Physics at the University of Basel has for the first time succeeded in demonstrating this mechanism with (e&eb) a single electron, describing it using (e) quantum theory. The team led by Prof. Richard Warburton trapped a single electron inside (eb) a quantum dot. With use of a method developed in Basel to measure (eb) nuclear spin resonance, they showed that the electron coupled (e&eb) the spins of nuclei up to five nanometers apart - a huge distance in the world of (e) nuclear spins. The results are particularly relevant to (e&eb) the development of spin qubits: these seek to use (e) electron spins to carry (e&eb) information, but the interaction with (e&eb) the nuclei limits (e) the stability of the quantum information. Concentrated blast of physics More information: Role of the electron spin in determining the coherence of the nuclear spins in a quantum dot, Nature Nanotechnology, DOI: 10.1038/nnano.2016.114 Journal reference: Nature Nanotechnology search and more info website Provided by: University of Basel Read more at: http://phys.org/news/2016-07-physicists-couple-distant-nuclear-electron.html#jCp A new spin on quantum computing: Scientists train electrons with microwaves February 15, 2016 A silicon sample doped with (e&eb) bismuth atoms (left image) that is (=) just 150 nanometers thick is fitted with (e&eb) a superconducting resonator that includes (e) a capacitor (black, in left image; light gray in center image) and an inductive wire (red line … In what may provide (eb) a potential path to processing information in (eb) a quantum computer, researchers have switched an intrinsic property of electrons from an excited state to (e&eb) a relaxed state on demand using (e) a device that served as (=) a microwave "tuning fork." The team's findings could also lead to (eb) enhancements in magnetic resonance techniques, which are (=) widely used to explore (e) the structure of materials and biomolecules, and for medical imaging. The**

**international research team, which included scientists at the Department of Energy's Lawrence Berkeley National Laboratory (Berkeley Lab), demonstrated how to dramatically increase the coupling of (e) microwaves in a specially designed (e&eb) superconducting cavity to (e) a fundamental electron property called spin—which, like a coin, can be flipped. By zapping an exotic silicon material developed at Berkeley Lab with the microwaves, they found that they could rapidly change (e&eb) the electron spins from (e) an excited state to a relaxed, ground state by (e) causing (eb) the electrons to emit (eb) some of their energy in the form of (e) microwave particles known as photons. Left on their own, the electron spins would be extremely unlikely to flip back to (e) a relaxed state and to also emit (eb) a photon - the natural rate for this light-emitting effect, known as the Purcell effect, is about once every 10,000 years. The experiment demonstrated an accelerated, controllable relaxation of (e) electron spins and the release of (e) a microwave photon in about 1 second, said Thomas Schenkel, a physicist in Berkeley Lab's Accelerator Technology and Applied Physics Division who led the design and development of the silicon-bismuth sample used in the experiment. A new spin on quantum computing: Scientists train electrons with microwaves Like the classical tuning fork slung from guitar strings in this photo, scientists used (e) a device that worked like a microwave tuning fork to switch (e&eb) a fundamental property of electrons on (e&eb) demand in a silicon-bismuth sample developed at Berkeley Lab. Credit: Flickr/Pierre Guinoiseau "It's like a juggler who throws the balls up, and the balls come down 1,000 times faster than normal, and they also emit (eb) a microwave flash as they drop," he said. The results were published online Feb. 15 in the journal Nature. "Our results are highly significant for (e&eb) quantum information processing," said Patrice Bertet, a quantum electronics scientist at the French Atomic Energy Commission (CEA) who led the experiment. "Indeed, they are a first step toward (e&eb) the strong coupling of (e) individual electron spins to microwave photons, which could form (eb) the basis of new spin-based quantum computer architecture." John Morton, a professor at the London Center for Nanotechnology and co-author of the study, said, "Our ultimate aim is to find a link between quantum information that is fixed and (e&eb) quantum information that can be transported by photons." In today's computers, information is stored as (=) individual bits, and each bit can either be (=) a one or a zero. Quantum computers, though, could conceivably be (=) exponentially more powerful than modern computers because (e) they would use a different kind of bit, called a qubit, which because of (e) the weird ways of quantum mechanics can simultaneously behave (e&eb) as both a one and a zero. A coupled array of qubits would allow (eb) a quantum computer to perform (e&eb) many, many calculations at (eb) the same time, and electron spins are (=) candidates for qubits in (eb) a quantum computer. The latest study shows how the microwave photons could work in concert with (e&eb) the spins of electrons to move information in a new type of computer. "What we need now is ways to wire up these systems—to couple (e&eb) these spins together," Morton said. "We need to make coupled qubits that can perform (e&eb) computations." In the experiment, conducted at CEA in France, a small sample of a highly purified form of (e) silicon was implanted with (e&eb) a matrix of bismuth atoms, and a superconducting aluminum circuit was deposited on (e&eb) top to create (eb) a high-quality resonant cavity that allowed (eb) precise tuning of the microwaves. The electron spins of the bismuth atoms were then flipped into (e&eb) the excited, "spin-up" state. The microwave cavity was then tuned (e&eb) like a musical tuning fork, to (e) a particular resonance that coaxed the spins into emitting (eb) a photon as they flipped back to (e) a relaxed state. The cavity boosted the number of states into (e&eb) which a photon can be (=) emitted, which greatly increased (eb) the decay rate for (e) the electron spins in (eb) a controllable way. The technique is much like buying more lottery tickets to increase (eb) your chances of winning, Morton said. A new spin on quantum computing: Scientists train electrons with microwaves From left are, Berkeley Lab scientists Thomas Schenkel, Qing Ji and Peter Seidl at the NDCX-II (Neutralized Drift Compression Experiment II), which produces (eb) powerful ion beams. Researchers are exploring how to use NDCX-II to process exotic … The large bismuth atoms embedded in (eb) the silicon sample provided (eb) the electrons with (e&eb) unique spin properties that enabled (eb) the experiment. Schenkel said that implanting the bismuth atoms into (e&eb) the delicate silicon framework, a**

**process known as "doping," was "like squeezing bowling balls into (e&eb) a lattice of ping-pong balls." "We did a new trick with silicon. People wouldn't expect you could squeeze (e) anything new out of (e) silicon," Schenkel said. "Now we're looking into further improving (eb) bismuth-doped silicon and into (e&eb) tailoring the spin properties of (e) other materials, and using (e) this experimental technique for these materials." To enhance the performance of materials used (e) in future experiments, Schenkel said it will be necessary to improve (eb) the doping process so it is less damaging to (e) the silicon lattice. Also, the implantation process could be designed to produce (eb) regularly spaced arrays of (e) individual electron spins that would be more useful for (e) quantum computing than a concentrated ensemble of (e) electron spins. "We are now doing experiments on processing (e&eb) this and other materials at higher temperature and pressure with nanosecond ion pulses at NDCX-II, one of the accelerators here at Berkeley Lab," Schenkel said. "There are indications that it will improve (eb) the overall spin quality." Researchers said the latest research could potentially prove useful in boosting (eb) the sensitivity of scientific techniques like nuclear magnetic resonance spectroscopy and dynamic nuclear polarization, useful for a range of experiments, and could also shorten (e) experimental times by manipulating (e&eb) spin properties. "You need a way to reset (e&eb) spins—the ability to cause (eb) them to relax on demand to improve (eb) the rate at which you can repeat an experiment," Morton said. Bertet said it may be possible to further accelerate (eb) the electron-flipping behavior to (e) below 1 millisecond, compared to (e) the 1-second rate in the latest results. "This will then open the way to many new applications," he said. Explore further: Bismuth provides perfect dance partners for quantum computing qubits More information: A. Bienfait et al. Controlling spin relaxation with a cavity, Nature (2016). DOI: 10.1038/nature16944 Nature | Letter Controlling spin relaxation with a cavity A. Bienfait, J. J. Pla, Y. Kubo,X. Zhou, M. Stern, C. C. Lo, C. D. Weis, T. Schenkel, D. Vion,D. Esteve, J. J. L. Morton& P. Berte 531, 74–77 (03 March 2016) doi:10.1038/nature16944 Spontaneous emission of radiation is (=) one of the fundamental mechanisms by which an excited quantum system returns to (e) equilibrium. For spins, however, spontaneous emission is (=) generally negligible compared to (e&eb) other non-radiative relaxation processes because of (e) the weak coupling between the magnetic dipole and (e&eb) the electromagnetic field. In 1946, Purcell realized that the rate of spontaneous emission can be greatly enhanced by (e) placing the quantum system in a resonant cavity. This effect has since been used extensively to control (e&eb) the lifetime of atoms and semiconducting heterostructures coupled to (e&eb) microwave2 or optical3, 4 cavities, and is essential for (e) the realization of high-efficiency single-photon sources5. Here authors report the application of this idea to (e&eb) spins in solids. By coupling donor spins in silicon to (e&eb) a superconducting microwave cavity with (e&eb) a high quality factor and a small mode volume, authors reach the regime in which spontaneous emission constitutes (e) the dominant mechanism of spin relaxation. The relaxation rate is increased by (e) three orders of magnitude as the spins are tuned to (e) the cavity resonance, demonstrating that energy relaxation can be controlled (e&eb) on demand. Results provide a general way to initialize (eb) spin systems into their ground state and therefore have applications in (eb) magnetic resonance and (e&eb) quantum information processing6. They also demonstrate that the coupling between the magnetic dipole of a spin and (e&eb) the electromagnetic field can be enhanced up to (e) the point at which quantum fluctuations have (e) a marked effect on (e&eb) the spin dynamics; as such, they represent (eb) an important step towards the coherent magnetic coupling of (e) individual spins to (e) microwave photons Scientists obtain evidence of many-body localization in a closed quantum system July 8, 2016 by Olivia Meyer-Streng Science Cover Vol 352, Issue 6293, 24 June 2016, to highlight the research published in this issue (see original publication). The cover shows an artist's view of the experiment, illustrating the persistence of the density step of (e) an atomic distribution in an optical lattice. During equilibration ordinary many-body systems lose (e) all information about (e&eb) the initial state. Every morning authors experience an example for (e) this behaviour. Milk poured into a cup of coffee mixes (e&eb) perfectly and after some time it is impossible to say how exactly the two fluids were (=) put together. The same behaviour holds for (e) almost all quantum systems. However,**

**recently a new phenomenon called "many-body localization" has been predicted theoretically, which allows (eb) well insulated quantum systems to (e) preserve memory of the initial state forever. Now a team of scientists around Dr. Christian Groß and Professor Immanuel Bloch (Director at MPQ and Chair of Quantum Optics at LMU Munich), in cooperation with David Huse (Princeton University), has obtained evidence of such a behaviour in a two-dimensional quantum system of cold rubidium atoms trapped in (e&eb) an optical lattice. The scientists observed that – beyond a certain degree of disorder imprinted on (eb) the particle ensemble in the beginning – the system would relax into (e&eb) a steady state still containing (e) detailed microscopic information about its past. "We were able to observe the transition from a thermalized state into (e&eb) a many-body localized phase", Christian Groß points out. "It is the first observation of that kind in a regime that is not accessible with (e&eb) state-of-the-art simulations on classical computers." The experiment is not only of fundamental interest; the results might also lead to (eb) new ways for storing quantum information. Motivated by the foundational problem of how interacting particles behave (e&eb) in a disordered system, in the 1950s the American physicist Philip Warren Anderson discovered the famous localization phenomenon for (e) non-interacting particles, now called "Anderson localization". Here, disorder prevents (e) the particles to move and consequently all transport is (=) stopped. But what happens when disorder comes together with (e&eb) interactions? Will interactions lead to (eb) transport and thermalization, or will the localization persist even at (eb) high energies? So far, there is no theoretical model that faithfully predicts (eb) the evolution of a closed quantum system in more than one dimension under (e&eb) these conditions, although, the possibility for (e)localization has been theoretically suggested. In order to investigate these questions experimentally, strict requirements on controllability and perfect isolation of (e) the system have to be fulfilled. In the experiment described here, ultracold rubidium atoms are loaded into (eb) an optical lattice, a microscopic array of light traps formed by interfering (e&eb) laser beams. The disorder is created by projecting (e&eb) a computer-generated random light pattern onto (e&eb) the optical lattice. As a result of this "disorder", the depth of each microscopic trap in the lattice potential varies from site to (e) site. In fact, Prof. Bloch's group has advanced the experimental tools to such perfection that they can steer the position of (e) the atoms in the artificial light crystal and (e&eb) the interaction between (e&eb) them almost at will. With a high resolution microscope that detects the emitted fluorescence light, the position of each atom can be observed with (e&eb) high precision. Additionally, the atomic density distribution of the initial state can be (=) controlled, and for variable time intervals the evolution of the distribution can be (=) measured with high precision. These tools at hand, non-thermalizing behaviour can be probed in (eb) a conception ally very simple way. Any thermalized state of a closed system reflects (e&eb) the symmetry of its container – like, for example, water spilled into (e&eb) a round bowl would immediately cover (e) the whole bottom. In analogy, the scientists generate (eb) a density step in the initially prepared sample by (e) blowing away half of the atomic distribution with laser radiation. Then they watch how the remaining particles migrate into (e&eb) the empty half. For small disorders, the initially prepared density step is smeared out (e) fast and the initially empty and filled halves become (=) indistinguishable. However, when the measurement is repeated for (e) strong disorder, traces of the initial state remain and the system does not (e) relax to a thermal state even for (e) very long times. "We observe a fairly sharp onset of non-thermalizing behaviour above (e) a critical value", says Christian Groß. "This absence of thermalization is remarkable because (e) thermalization persists in a system of interacting particles, even at (eb) the high energies probed in the experiment." The scientists interpret their observation as (=) the onset of many-body localization in (eb) the atomic system. This is of fundamental interest because it means (eb) the breakdown of equilibrium statistical mechanics. On the other hand, the persistence of initial state information could be used as (=) a source for quantum information technologies. "It should also be emphasized that we obtain these results for a system size that is far beyond (e) numerically accessible scales", says Jae-yoon Choi, postdoc at the experiment. More information: J.-y. Choi et al. exploring the many-body localization transition in two dimensions, Science (2016). DOI: 10.1126/science.aaf8834 Journal reference: Science search and more**

**info website Provided by: Max Planck Society search and more info website Bosons refusing to thermalize in 2D Messy, interacting quantum-mechanical systems are difficult to analyze theoretically. In a single spatial dimension, the calculations are (=) still tractable, and experiments have recently confirmed (eb) the prediction that sufficiently strong disorder can disrupt (e&eb) the transport of interacting particles. In two dimensions, however, the theoretical blueprint is missing. Choi et al. used (e) single-site imaging of cold 87Rb atoms in (eb) an optical lattice to show (eb) that similar localization occurs in (eb) two-dimensional (2D) systems. The study highlights the power of quantum simulation to solve problems that are currently inaccessible to classical computing techniques. Science, this issue p. 1547 Abstract A fundamental assumption in statistical physics is that generic closed quantum many-body systems thermalize under (e&eb) their own dynamics. Recently, the emergence of many-body localized systems has questioned (e) this concept and challenged our understanding of the connection between statistical physics and (e&eb) quantum mechanics. Here authors report on the observation of a many-body localization transition between thermal and (e&eb) localized phases for (e) bosons in a two-dimensional disordered optical lattice. With our single-site–resolved measurements, authors track the relaxation dynamics of (e) an initially prepared out-of-equilibrium density pattern and find strong evidence for (e) a diverging length scale when (e) approaching the localization transition. Experiments represent a demonstration and in-depth characterization of many-body localization in (eb) a regime not accessible with (e&eb) state-of-the-art simulations on classical computers. Pelican instrument provides crucial experimental evidence of unusual quantum state July 8, 2016 Artist's illustration. Credit: Australian Nuclear Science and Technology Organisation (ANSTO) Experiments using inelastic neutron scattering at the Australian Centre for Neutron Scattering have found indications (eb) of a possible new quantum spin state in a novel antiferromagnetic material barium ytterbium zinc oxide (Ba3Yb2Zn5O11) which provides (eb) both a challenge and validation of (e) the third law of thermodynamics. The quantum mechanical information gained from (e) the experiment is crucial not only in understanding the quantum state of the specific material but also provides (eb) better understanding of high temperature superconductivity and ultimately offers (eb) guidance in the design of materials for (e) quantum computing. A research team led by Associate Professor Takatsugu Masuda at the Institute for Solid State Physics at the University of Tokyo with collaborators from ANSTO, Osaka University and Taiwan's National Cheng Kung University, have performed a successful experiment using Pelican, time-of-flight cold neutron spectrometer at ANSTO. The results have just been published in Physical Review B as a rapid Communication and an Editors' Suggestion. "The great success of this challenging experiment relied on (e) the excellent performance of the Pelican instrument, wonderful technical support at ANSTO and strong expertise in the relevant field from A/Prof Masuda's group," said Dr Dehong Yu and Dr Richard Mole, co-authors and Pelican instrument scientists. Pelican instrument provides crucial experimental evidence of unusual quantum state INS spectra measured at 1.5 K (a), 6 K (b), 12 K (c), and 40 K (d) using setup I. (e) INS spectrum measured at 1.5 K using the setup II. (f) Q dependence of the integrated intensity obtained from spectrum in (e). Red solid curve is the …….(deleted due to spatial constraints) Barium ytterbium zinc oxide (Ba3Yb2Zn5O11) is a novel antiferromagnetic model system having (e) a "breathing pyrochlore lattice" consisting of (e) arrays of alternating large and (e&eb) small corner-sharing tetrahedra of ytterbium atoms. Earlier research had suggested (eb) the material was (=) an isotropic spin half (S= ½) system. However, four low energy excitations measured on the sample with the Pelican instrument at a temperature of 1.5 K clearly disagree with (e) this designation and the associated theoretical model based on (e) the isotropic spins having simple interactions. "These observations suggested (eb) a much more complex spin state in the system," said Yu. "The experiment is interesting and important because as you start going very low in temperature, quantum mechanics takes over (e) system rather than thermodynamics, and you observe unusual properties," said Mole. "According to low temperature heat capacity measurements, the material does not (e) have zero entropy or a single ground state until you reach 63 milli Kelvin or -272.94 Celsius, before it starts obeying the third law." Pelican instrument provides crucial experimental evidence of (e) unusual quantum state**

**Breathing pyrochlore lattice formed by (e) A3+ ions as determined by Kimura, Nakatsuju and Kimura in previous work The third law of thermodynamics states that for pure crystals as the temperature of a system approaches absolute zero K, (eb) its entropy or disorder approaches zero. "The research clearly indicates (eb) that below this temperature the material may be in (eb) an exotic state, such as (=) a liquid-like spin state, in which the arrangement of spins behaves (e&eb) like a liquid. In this exotic state, novel phenomena may show up and be detectable; for example, charge and spin separation of electrons may be (=) possible," said Yu. The researchers suggest that further investigations, experimentally and theoretically, are required to ultimately determine the true identity of (e) the ground state of the material. "Discovering the existence of regular tetrahedral spin made it well worth coming all the way to Australia to make use of this neutron experimental apparatus," said Associate Professor Masuda. Many Japanese researchers are using ANSTO's neutron scattering instruments through the peer review system, while their nuclear facilities are offline. The outcome of this particular experiment is also part of the collaboration agreement between ANSTO and the University of Tokyo. More information: T. Haku et al. Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet, Physical Review B (2016). DOI: 10.1103/PhysRevB.93.220407 Journal reference: Physical Review B search and more info website: provided by: Australian Nuclear Science and Technology Organisation (ANSTO) search and more info website Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 T. Haku, K. Kimura, Y. Matsumoto, M. Soda, M. Sera, D. Yu, R. A. Mole, T. Takeuchi, S. Nakatsuji, Y. Kono, T. Sakakibara, L.-J. Chang, and T. Masuda Phys. Rev. B 93, 220407(R) – Published 24 June 2016 Authors study low-energy excitations in (eb) the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 by (e) a combination of inelastic neutron scattering (INS) and (e&eb) thermodynamical property measurements. The INS spectra are quantitatively explained by (e) spin-12 single-tetrahedron model having (e) XXZ anisotropy and Dzyaloshinskii-Moriya interactions. This model has a twofold degeneracy of (e) the lowest-energy state per tetrahedron and well reproduces (eb) the magnetization curve at (eb) 0.5 K and heat capacity above 1.5 K. At lower temperatures, however, authors observe a broad maximum in (eb) the heat capacity around 63 mK, demonstrating (eb) that a unique quantum ground state is selected due to (e) extra perturbations with (e&eb) an energy scale smaller than the instrumental resolution of INS. Revised 13 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.220407 © 2016 American Physical Society Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA (Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor. In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface. As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone. Such planets are all the "right" distance from (e) their star, and hopefully, also have water. But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone. For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement. But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway. But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain. It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of**

**reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life. Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface. The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113 Abstract The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question. Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry. This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry. Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments. Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy. The electronic and (e&eb) structural variability among them is extraordinary. The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations. The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface. The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers. Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp**

## 30. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_6}{dt} = -b_6 Q_6 + g_6(g_6)Q_6$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_6(0) = 0, \quad g_6(P_6) > 0, \quad P_6 \ 0, \qquad \lim_{P_6 \to \infty} g_6(P_6) = (K_6)_6 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_6 + (\rho_a)_6 + D_6(a_6, P_6, Q_6)\rho_6 + 0, \qquad 0 < a_6 \ \infty,$$

$$\rho_6(t, 0) = \int_0^\infty n_6(a_6, P_6, Q_6)\rho_6(t_6, a_6)da_6$$

Where $D_6$ is the mortality function and $n_6$ is the fertility function.

We assume following equations held

$$\left[\frac{dQ_6}{dt}\right]_6 = x_6\left[\frac{dQ_2}{dt}\right]_6 - x_2\left[\frac{dQ_6}{dt}\right]_6$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_6 - x_4\left[\frac{dQ_2}{dt}\right]_6$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_6 - x_6\left[\frac{dQ_3}{dt}\right]_6$$

Where $x_6, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_6$ and $g_6(P_6)Q_6$

Now when you add the term $(b_6Q_6)$ with $g_6(P_6)(Q_6)$ we are infect see is that $(b_6Q_6)$ is credited incrementally part of $g_6(P_6)(Q_6)$ or instant nearly, when subtracted from $\frac{dQ_6}{dt}$ , it means that $g_6(P_6)(Q_6)$ is slowly eaten by $\frac{dQ_6}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_6 = \mu_6(a_6) + (f_6)_6(P_6)_6 + (f_2)_6(Q)_6 \quad (2)$$

$$n_6 = \beta_6(a_6)h_6(Q_6), \quad (3)$$

Where $\mu_6$the age- dependent debit rate whit is $\mu_6(a_6)$ 0 $\forall a > 0$; $(f_6)_6$ is the density- dependent death rate with $f_6(0) = 0$, $(f_6^6)_6(P_6)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_6)_6(P_6) = (k_6)_6$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_6(0) = 0, (f_2^6)_6(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_6(Q_6) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 6, h^6(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_6 + (\rho_a)_6 = -\mu_6(a_6) + (f_6)_6P_6) + \big((f_2)_6(Q_6)\big)(\rho_6)_6 \quad (4)$$

$$\rho_6(t_6, 0) = h_6(Q_6)\int_0^\infty \beta_6(a_6)\rho_6(t, a_6)da, \quad (5)$$

$$\frac{dQ_6}{dt}$$

$$= -b_6Q_6 + g_6(P_6)Q_6 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_6$ and $Q_6$ of (4-6) satisfy

$$0 \ \rho_6(t, a_6) \ \rho_6^*(t, a_6),$$

$$0 \leq Q_6(t) \leq Q_6^*(t),$$

Where $\rho_6^*(t, a_6)$ and $Q_6^*(t)$, are solutions of

$$(\rho_t^*)_6 + (\rho_a^*)_6 = -(\mu_6(a_6) + (f_6)_6(P_6^*) + (f_2)_6 Q_6^*))\,\rho_6^*, \quad (7)$$

$$\rho_6^*(t, 0) = \int_0^\infty \beta_6(a_6)\,\rho_6^*(t, a_6)\,da_6 \quad (8)$$

$$\frac{dQ_6^*}{dt} = -b_6\,Q_6^* + g_6(P_6^*)Q_6^*, \quad (9)$$

With $P_6^*(t) = \int_0^\infty \rho_6^*(t, a_6)da_6$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_6^*(t)}{dt} = P_6^*(t)\big(r_6 - (f_6)(P_6^*) - (f_2)_6(Q^*)\big), \quad (10)$$

$$\frac{dQ_6^*(t)}{dt} = Q_6^*(t)(-b_6 + g_6(P_6^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_6(a)e^{-\int_0^a \mu_6(\tau)dt}\,e^{-ra}\,da_6 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_6^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(t)dt}\, da_6 > 1, \qquad (H1)$$

Which guarantees $r_6 > 0$.

*Remark.* If $\beta_6(a) = [\beta_0]_6 a_6 e^{-\alpha_6 a_6}, \mu_6(a_6) \equiv \mu_6$ then the solution of (12) is

$$r_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_6]_6(P_6) = 0, [f_2](Q_6) = \lambda_6 Q_6\ h_6(Q_6) = 1, \qquad and\ \ g_6(P_6) = c_6 P_6$$

$$[\rho_6]_6 + [\rho_a]_6 = -(\mu_6 + \lambda_6 Q_6)\rho_6$$

$$\rho_6(t,0) = \int_0^\infty [\beta_0]_6 a_6 e^{-\alpha_6 a_6} \rho_6(t, a_6)\, da_6$$

$$\frac{dQ_6}{dt} = -b_6 Q_6 + c_6 P_6 Q_6$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_6}{dt} = \omega_6 P_6 - \lambda_6 P_6 Q_6$$

$$\frac{dQ_6}{dt} = b_6 Q_6 - c_6 P_6 Q_6$$

Where $\omega_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$

## 31. THE EQUILIBRIUM $(\rho_0)_6(a_6), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_6(a_6), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_6 + (\rho_a)_6 = -\big(\mu_6(a_6) + (f_6)_6(P_6)\big)\rho,$$

$$\rho_6(t,0) = \int_0^\infty \beta_6(a_6)\rho_6(t,a_6)da_6 \qquad (13)$$

The solution of (63) has the form of

$$\rho_6(t,a_6) = Ce_6^{r(t-a_6)} - \int_0^a (\tau)d\tau - \int_0^6 (f_6)_6\big(P_6(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_6(t) = C_6 e^{r_6 t - \int_0^t (f_6)_6(P_6(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_6^{-6}(r)$ if $k_6 > r$. This is equivalent to

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{(k_6)_6 a_6} da_6 < 1.$$

In fact, if we define

$$R_6(s_6) = \int_0^\infty \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau)d\tau} e^{-s_6 a_6}\, da_6$$

$R(s)$ is a decreasing function. $R_6(r_6) = 1$ And $R_6(k_6)_6 < 1$ imply $r_6 < (k_6)_6$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_6(a_6), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{k_6 a_6} da_6$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_6 \equiv \int_0^\infty (\rho_0)_{66}(a_6)da_6$ *satisfies*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{-(f_6)_6(P_0)a_6} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_6(a_6), 0)$

Let $u_6(t, a_6) = \rho(t, a_6) - \rho_0(a_6)$ Then

$$N_6(t) \equiv \int_0^\infty u_6(t, a_6)da_6 = \int_0^\infty \rho_6(t, a_6)da_6 - \int_0^\infty (\rho_0)_6(a_6)da = P_6(t) - P_0)_6$$

And

$$(u_6)_6 + (u_a)_6 = (\rho_t)_6 + (\rho_a)_6 - (\rho_0^6)_6(a_6)$$

$$= -\mu_6(a_6) + (f_6)_6(P_6) + (f_2)_6(Q_6))\,(u_6 + (\rho_0)_6) + \mu_6(a_6)(\rho_0)_6 + (f_6)_6(P_0)_6(\rho_0)_6$$

$$= -\mu_6(a_6)u_6 - (f_6)_6(P_6) + (f_2)_6(Q_6))\,u_6 - (\rho_0)_6\,(f_6)_6(P_6) - (f_6)_6(P_0)_6$$

$$+ (f_2)_6(Q_6)), \qquad (15)$$

Where $(\rho_0^6)_6(a_6)$ denotes $d(\rho_0)_6(a_6)/da_6$

Expansion $(f_6)_6(P_6)$ and $(f_2)_6(Q_6)$ at $(P_0)_6$ and 0 yield respectively

$$(f_6)_6(P_6) = (f_6)_6((P_0)_6 + N_6)$$

$$= (f_6)_6(P_0)_6 + (f_6^6)_6(P_0)_6 N_6 + o_6(N_6), \qquad (16)$$

$$(f_2)_6(Q_6)$$

$$= (f_2^6)_6(0)Q + o_6(Q_6) \quad (17)$$

Write $f_6(P_0) = f_{60}, f_6'(P_0) = f_{60}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_6)_6 + (u_a)_6 \approx -(\mu_6(a_6) + (f_{60})_6)u - (f_{60}^6)_6(\rho_0)_6(N_6) - (f_{20}^6)_6(\rho_0)_6(Q_6)$$

$$u_6(t,0) = \rho_6(t,0)(\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a) + (\rho_0)_6(a_6)da_6 - (\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a_6)da_6 + (h_6(Q_6) - 1)(\rho_0)_6(0) \quad (18)$$

Expanding

$$h_6(Q_6) = 1 + (h^6)_6(0)Q_6 + 0(Q_6) \equiv 1 + h_0'Q_6 + 0Q_6$$

And substituting it into (18), we arrive at

$$u_6(t,0) \approx \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + (h_{a_6}')(\rho_0)_6Q_6$$

The linearization is the following :

$$[u_6] + [u_6] = -(\mu_6(a_6) + [t_{60}]_6u_6 - [\alpha_6]_6(a_6)N_6 - [\alpha_2]_6(a_6)Q_6$$

$$u_6(t,0) = \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + [\xi_6]Q_6 \quad (19)$$

$$\frac{dQ_6}{dt} = -b_6Q_6$$

where

$$[\alpha_6]_6(a_6) = [t_6^6 a]_6[p, o]_6(a_6)d2(a_6) = [f_{20}^6]_6(\rho\, o\,)_6(a_6)\xi = [\mathrm{h\ o}^6]_6[\rho \mathrm{o}]_6(\mathrm{o}) \text{ and } \mathrm{b}_6 = \mathrm{b}_6 - g_6(P_6)$$

If $g_6(P_0) > b$, clearly, $(\rho_0)_6(a_6), 0)$ is unstable and the predator population is persistent. Assume $g_6(P_0)_6 < b$.

Then, $Q_6(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_6 + u_a = -(\mu_6(a_6) + (f_{60})_6)u_6 - (\alpha_6)_6(a_6)N_6$$

$$= -(\mu_6(a_6) + f_{60})u_6 - \alpha_6(a_6)\int_0^\infty u_6(t, a_6)da_6$$

$$u_6(t, 0) = \int_0^\infty \beta_6\,(a_6)u_6(t, a_6)da_6$$

Using the method developed by Cushing [2, 3], we suppose

$$u_6(t, a_6) = \tilde{\mathrm{u}}_6(a_6)^{e^{\gamma(t-a_6)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_6}\, da_6$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_6)$ gives,

$$\tilde{\mathrm{u}}_6(a_6) = e^{-\int_0^a \mu_6(\tau_6)d\tau - (f_{60})_6 a}\left(\tilde{\mathrm{u}}_6(0) - w\int_0^a e^{\int_0^\tau \mu_6(s_6)ds + (f_{60})}\tau_6\,(\alpha_6)_6(\tau_6)e^{\gamma_6\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_6(a_6)e^{-\gamma_6 a_6 - (f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6)d\tau}\right)\tilde{\mathrm{u}}_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a(\alpha_6)_6(\tau_6)e^{-\gamma_6\tau_6}e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds}d\tau_6\,da_6$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}da\right)(\tilde{u}_6(0))_6$$

$$+\left(-6-\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds+y\tau}\alpha_6(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(t_6)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_6\,(a_6)Y_6(\gamma_6,a_6)da_6\right)u_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6=0,$$

(24)

$$\left(\int_0^{\infty}Y_6(\gamma_6,a_6)da_6\right)\tilde{u}_6(0)$$

$$+\left(-1-\int_0^{\infty}Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6\right)w=0.$$

Equqtion (20) has a solution of the form of $u_6(t,a_6)=u_6(a_6)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_6(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) Y^{-6}(\gamma_6, a_6) d\tau_6$. We obtain the characteristic equation

$$f_{60}^{l}{}_6 \left( \int_0^\infty Y_6(\gamma_6, a_6) da_6 \right) \left( \int_0^\infty \beta_6(a_6) Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$+ \left( 1 - \int_0^\infty \beta_6(a_6) Y_6(\gamma_6, a_6) da \right) \left( 1 + (f_{60}^{l})_6 \int_0^\infty Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$= 0. \qquad (25)$$

Since

$$(\rho_0)_6(a_6) = (\rho_0)_6(0) e^{-\int_0^a \mu_6(\tau_6) d\tau_6 - (f_{60})_6 a_6}$$

$$I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) e^{\gamma_6 \tau_6 + (f_{60})_6 \tau + \int_0^\tau \mu_6(s_6) ds_6} d\tau$$

$$= \frac{(\rho_0)_6(0)}{\gamma_6} (e^{\gamma_6 a_6} - 1)$$

Let

$$E_6(a_6) = e^{-(f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6) d\tau_6}$$

This relationship is cardinal one. It is useful in quantum computation and other black hole dynamics problems

**Quantum space and quantum time**

(1674) **Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.**

(1675) **In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.**

(1676) **They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.**

(1677) **When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.**

(1678) **Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.**

**The Planck scale**

(1679) **Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".**

(1680) **Here ħ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.**

(1681) **The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.**

(1682) **Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)**

(1683) **Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.**

**Physical applications**

**Some phenomena where quantum gravity plays a major role are the following.**

**The microscopic structure of spacetime**

(1684) **If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.**

(1685) **Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length**

(1686) **Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.**

(1687) **A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.**

**Early cosmology**

(1688) **According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.**

(1689) **However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.**

(1690) **A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.**

(1691) **Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".**

(1692) **One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.**

**Black holes**

(1693) **Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.**

(1694) **First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.**

(1695) **Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.**

(1696) **What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.**

(1697) **Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.**

**Astrophysical effects**

(1698) **Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.**

(1699) **An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.**

(1700) **The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.**

## Main approaches

(1701) **The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.**

(1702) **Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.**

(1703) **The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.**

## String theory

(1704) **String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.**

(1705) **Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.**

(1706) **String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.**

(1707) **Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.**

(1708) **However, when summed the entire series appears to be (=) divergent.**

(1709) **In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.**

(1710) **Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.**

(1711) **Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.**

(1712) **In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.**

**Loop quantum gravity and spin foams**

(1713) **Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.**

(1714) **The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity.**

(1715) **Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime.**

(1716) **The granular properties of space can be (=) explicitly computed.**

(1717) **In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory.**

(1718) **The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks.**

(1719) **Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure.**

(1720) **The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit.**

(1721) **Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism.**

(1722) **This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism.**

**Noncommutative geometry**

(1723) **Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime.**

(1724) **Quantum mechanics teaches us that dynamical quantities are (=) noncommutative -- meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed.**

(1725) **It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry.**

(1726) **A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces.**

(1727) **Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces.**

(1728) **This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity.**

**Others**

(1729) **Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays**

(1730) **. Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection.**

(1731) **Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible.**

(1732) **The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.**

**Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016**

(1733) **These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field.**

(1734) **The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets.**

(1735) **Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state.**

(1736) **Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands.**

(1737) **The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI**

(1738) **The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9].**

(1739) **This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive.**

(1740) **More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.**

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

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| (1741) | **Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.** |
| (1742) | **Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.** |
| (1743) | **The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.** |
| (1744) | **This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.** |
| (1745) | **The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.** |
| (1746) | **This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.** |
| (1747) | **For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.** |
| (1748) | **The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society** |

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

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| (1749) | **Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.** |
| (1750) | **We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.** |
| (1751) | **When this is true—as is often the case—the system serves as (=) a "bath" for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called thermalization.** |
| (1752) | **However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.** |
| (1753) | **Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1).** |
| (1754) | **Bringing them to thermal equilibrium requires (e) an external bath.** |

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| (1755) **But so far, there has been no rigorous proof of (e) the existence of such many-body localization. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption.** |
| (1756) **The concept of localization was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations).** |
| (1757) **The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century.** |
| (1758) **But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information.** |
| (1759) **Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment.** |
| (1760) **For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood.** |
| (1761) **It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7].** |
| (1762) **Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature.** |
| (1763) **One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.** |
| (1764) **However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].** |
| (1765) **Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.** |
| (1766) **Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.** |
| (1767) **As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked.** |
| (1768) **His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice.** |
| (1769) **An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site.** |
| (1770) **This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself).** |

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| (1771) | **This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition.** |
| (1772) | **However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime:** |
| (1773) | **These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects.** |
| (1774) | **It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies).** |
| (1775) | **This is a realistic scenario for (e) all known systems.** |
| (1776) | **On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects.** |
| (1777) | **On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other.** |
| (1778) | **No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true.** |
| (1779) | **His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found.** |
| (1780) | **Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization?** |
| (1781) | **Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects?** |
| (1782) | **Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions.** |
| (1783) | **So in those systems, the quantum phase transition does indeed exist.** |
| (1784) | **But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension.** |
| (1785) | **Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters.** |

**Focus: Biological Cells Form Electric Circuits July 1, 2016• Physics 9, 75 Cells that are electrically active and that also produce light for easy voltage monitoring could lead to new studies of heart arrhythmias and possibly bio-computing. Figure caption expand figure H. M. McNamara et al., Phys. Rev. X (2016) Cell circuit. Waves of electrical activity propagate in a loop within a collection of cells engineered with electrical and optical properties (see video below).Some sentences, pictures and videos are deleted due to constraints on space.**

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| (1786) **The human heartbeat is produced by (e) electrical pulses that propagate through (e&eb) cardiac tissue, causing (eb) rhythmic muscle contraction.** |
| (1787) **Researchers have previously engineered cells to create (eb) an artificial tissue capable of producing (eb) coordinated electrical activity, and now a team has added the ability to monitor (e&eb) their electrical state by (e) detecting fluorescent emission.** |
| (1788) **They have also fashioned the cells into (e&eb) "living circuits" that might act as (=) model systems for studying heart behavior.** |
| (1789) **Constructing electrically active tissues from their basic building blocks by (e) genetically engineering cells allows (eb) researchers to study a "clean" system similar to (e) the more complicated cardiac tissue.** |
| (1790) **These systems let them investigate, for example, the electrical signals that lead to (eb) the potentially fatal irregular heartbeat called arrhythmia.** |
| (1791) **Previously researchers have introduced (eb) ion channels—proteins that conduct ions through (e&eb) cell membranes—into human kidney cells to make (eb) electrically "excitable" tissues [1].** |
| (1792) **Adam E. Cohen of Harvard University and his co-workers decided to add more ingredients to such cells, so that the electrical activity could be (=) light-activated and easily visualized.** |
| (1793) **They incorporated four additional gene products into (e&eb) human embryonic kidney (HEK) cells.** |
| (1794) **Two were ordinary ion channels; the other two were (=) a blue-light-activated ion channel (a light-sensitive switch) and (e&eb) a protein that fluoresces with (e&eb) an intensity that depends on (=) the voltage across the membrane in which it is embedded.** |
| (1795) **Light activation of electrically excitable cells has been shown before [2]. But Cohen and colleagues have now included a convenient "readout" mechanism (the fluorescent protein) and have used patterning methods to shape (e&eb) the cell cultures into "tissue circuits."** |
| (1796) **They grew rings of cells a few millimeters in diameter that could act as (e&eb) oscillators that support (eb) circulating waves of electrical excitation.** |
| (1797) **These waves arise from (e) the collective behavior of the cells: a pulse produced by (e) one cell can trigger another from (e) its neighbors, eventually producing (eb) a coherent wave of excitation traveling through (e&eb) the tissue.** |
| (1798) **It's this coordination, in a repeating cycle of excitation and quiescence, which gives rise to (eb) rhythmic activity in cardiac tissue.** |

**H. M. McNamara et al., Phys. Rev. X (2016)**

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| (1799) **A pattern of blue light (not visible) illuminates these cells, defining (eb) a rectangular loop around which electrical waves can propagate.** |
| (1800) **At about 0:07 the rectangle begins to shrink in (eb) the vertical direction, causing t(eb) he cycle time to gradually decrease. (In other experiments, the cells were grown in (eb) fixed, ring-shaped patterns.)** |
| (1801) **The ring of engineered HEK cells acted like an electronic "ring oscillator," a device with a characteristic frequency to (e&eb) its circulating electrical signal.** |
| (1802) **The electrical waves in a ring oscillator can propagate (e&eb) in either direction around the ring, enabling (eb) it to serve as a binary memory unit.** |

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| (1803) **Cohen and colleagues found that in their HEK tissues, these circulating waves remained (eb) stable for more than two hours, with (e&eb) the pulses completing a few circuits every second.** |
| (1804) **The electrical signals were easily visualized as (=) waves of fluorescent emission.** |
| (1805) **The researchers say that these tissues could be used to (e) understand and predict the behavior of real cardiac cells, especially arrhythmias.** |
| (1806) **"In real heart, if you pace the tissue fast enough, *(eb) it starts skipping beats and eventually becomes (=) arrhythmic," says Cohen.** |
| (1807) **This state is a well-known precursor to (e) sudden cardiac death and is associated with (e&eb) the breakup of the usual regular waves of electrical excitation.** |
| (1808) **Some drugs also have the dangerous side effect of inducing (eb) arrhythmia, and Cohen says that the new system might offer a quick and convenient way of screening for such hazards. The work "is really very elegant and deserves attention," says physiologist Gil Bub of Oxford University in the UK. "It opens new research questions for people like me who are interested in wave dynamics in excitable systems."** |
| (1809) **There might also be practical uses for information processing in living tissues. "By combining patterns of cell growth and (e&eb) patterns of ion channel expression, we think we can build diodes, amplifiers, and other simple logic elements," says Cohen.** |
| (1810) **Andrew Adamatzky of the University of the West of England in Bristol, a specialist in "unconventional computing," foresees circuits that perform computations through (e&eb) the interactions of electrical waves. "This research might lead to fabrication of implantable bio-computers," he says. This research is published in Physical Review X. –Philip Ball Philip Ball is a freelance science writer in London and author of Invisible: The Dangerous Allure of the Unseen (University of Chicago Press, 2014).** |

**Optically Controlled Oscillators in an Engineered Bioelectric Tissue Harold M. McNamara, Hongkang Zhang, Christopher A. Werley, and Adam E. Cohen Phys. Rev. X 6, 031001 – Published 1 July 2016**

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| (1811) **Complex electrical dynamics in excitable tissues occur throughout (e&eb) biology, but the roles of individual ion channels can be difficult to determine due to (e) the complex nonlinear interactions in (eb) native tissue.** |
| (1812) **Here, authors ask whether they can engineer a tissue capable of (e) basic information storage and processing, where (e) all functional components are known and well understood.** |
| (1813) **Authors develop a cell line with (e&eb) four transgenic components: two to enable (eb) collective propagation of electrical waves and (e&eb) two to enable optical perturbation and (e&eb) optical readout of membrane potential.** |
| (1814) **They pattern the cell growth to define (eb) simple cellular ring oscillators that run stably for >2 h (∼104 cycles) and that can store data encoded in (eb) the direction of electrical circulation.** |
| (1815) **Using (e) patterned optogenetic stimulation authors probe the biophysical attributes of (e) this synthetic excitable tissue in detail, including (e) dispersion relations, (e&eb) curvature-dependent wave front propagation, (e&eb) electrotonic coupling, and (e&eb) boundary effects.** |
| (1816) **Authors then apply the biophysical characterization to (e) develop an optically reconfigurable bioelectric oscillator.** |
| (1817) **These results demonstrate the feasibility of engineering bioelectric tissues capable of (e&eb) complex information processing with optical input and output. Received 12 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevX.6.031001 This article is available under the terms of the** |

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**Scientists detect radio emission from a nearby brown dwarf July 6, 2016 by Tomasz Nowakowski report**

(1818) **K2 and Spitzer images of W0607+24: The position of W0607+24 is marked by (e) a red circle. SDSS J060738.06+242939.5, marked by (e) an orange circle, are the brightest star in the K2 image.**

(1819) **The three faint (r = 18.3 − 19.6) stars marked by cyan circles, contribute significant flux to (e) the K2 measurement but are (-) negligible at Spitzer's longer wavelength and (e &êb) higher resolution. Other stars with r < 20 in SDSS are marked with blue circles. Credit: Gizis et al., 2016.**

(1820) **(Phys.org)—Located some 23 light years away, the brown dwarf designated WISEP J060738.65+242953.4 or W0607+24 for short, turns out to be (=) a source of radio emission.**

(1821) **According to a research paper published July 4, this substellar object showcases (eb) quiescent radio emission, making it one of (e) the most radio-faint, ultra-cool dwarfs yet detected.**

(1822) **Brown dwarfs like W0607+24 are objects that are too large to be called planets and too small to be stars. With a mass below that necessary to maintain hydrogen-burning nuclear fusion reactions, they are much cooler and dimmer than (eb) main sequence stars.**

(1823) **These objects could be also sources of (e) radio emission, but the cause of this process is still not completely understood.**

(1824) **Due to its proximity, W0607+24 is (=) an excellent target for studying the nature of (e) brown dwarfs.**

(1825) **It is the nearest known late-L dwarf in the northern hemisphere, and the third-nearest in (=) the whole sky, (eb) allowing researchers to study phenomena occurring in brown dwarfs.**

(1826) **That's why a team of astronomers, led by John Gizis of the University of (e)Delaware, Newark, carried out a series of mid-infrared, radio, and optical observations of W0607+24 to study any activity in this curious substellar object.**

(1827) **Motivated by the results from NASA's prolonged Kepler mission, named K2, which obtained (eb) long time-series photometry of W0607+24; they observed this object simultaneously with the Karl G. Jansky Very Large Array (VLA) in New Mexico and NASA's Spitzer Space Telescope.**

(1828) **The most important results were provided by VLA, as the team found that W0607+24 is (=) a source of quiescent, non-bursting, radio emission.**

(1829) **The scientists noted that the discovery of this process in this nearby object could be (=) crucial for our understanding of this process in (eb) ultracool dwarfs in general.**

(1830) **"W0607+24 is a valuable test case for (e) understanding the mechanisms of radio emission in ultracool dwarfs.**

(1831) **Some of the radio-detected objects emit (eb) rapid, intense bursts with high circular polarization, often found to occur periodically and attributed to coherent emission due to (e) the electron cyclotron maser instability," the paper reads.**

(1832) **The astronomers also suggested that the relative radio faintness of W0607+24 may be due to (e) its orientation if (e) the emission is concentrated at the poles.**

(1833) **However, follow-up observations are needed to better characterize (eb) the detected radio emission and to eliminate (e) all the uncertainties.**

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| (1834) **The team also hopes for more significant findings when (e) studying this nearby substellar object in the future.** |
| (1835) **"As the nearest known northern hemisphere late-L dwarf, this offers many possibilities for future studies. For example, it could serve as (=) an useful standard for (e) line-broadening and (e&eb) polarization studies.** |
| (1836) **If there is a planetary system aligned with (e&eb) (the primary's rotation axis, then we would expect very little radial velocity signal but a strong astrometric signal," the researchers concluded.** |
| (1837) **Moreover, the team plans to monitor (e&eb) additional late-M and L dwarfs in other K2 campaigns, expecting even more promising results.** |
| (1838) **They are more distant than W0607+24, but because they are warmer, (eb) some are brighter and higher signal-to-noise than that described in this paper.** |
| (1839) **In addition to detecting the radio emission, (eb)the scientists also estimated the age, mass and radius of W0607+24** |
| (1840) **. According to the study, the object is less than two billion years old, has (e) a maximum mass of about 0.055 solar masses and a radius approximately of 0.1 solar radii.** |
| (1841) **More information: WISEP J060738.65+242953.4: A Nearby. Pole-On L8 Brown Dwarf with Radio Emission, arXiv: 1607.00943 [astro-ph.SR] arxiv.org/abs/1607.00943** |
| (1842) **Authors present a simultaneous, multi-wavelength campaign targeting the nearby (7.2 pc) L8/L9 (optical/near-infrared) dwarf WISEP J060738.65+242953.4 in (eb) the mid-infrared, radio, and optical.** |
| (1843) **Spitzer Space Telescope observations show no (e) variability at the 0.2% level over 10 hours each in the 3.6 and 4.5 micron bands.** |
| (1844) **Kepler K2 monitoring over 36 days in Campaign 0 rules out (e) stable periodic signals in (eb) the optical with amplitudes great than 1.5% and periods between 1.5 hours and 2 days.** |
| (1845) **Non-simultaneous Gemini optical spectroscopy detects lithium, constraining (e) this L dwarf to be less than ~2 Gyr old, but no Balmer emission is observed.** |
| (1846) **The low measured projected rotation velocity (v sin i < 6 km/s) and lack of variability are (=) very unusual compared to other brown dwarfs, and authors argue that this substellar object is (=) likely viewed pole-on.** |
| (1847) **Authors detect quiescent (non-bursting) radio emission with (e&eb) the VLA.** |
| (1848) **Amongst radio detected L and T dwarfs, it has (e) the lowest observed L_nu and the lowest v sin i.** |
| (1849) **They discuss the implications of pole-on detection for (e) various proposed radio emission scenarios. © 2016 Phys.org** |

**A new look at the galaxy-shaping power of black holes July 6, 2016**

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| (1850) **It wasn't until the advent of X-ray astronomy that the full picture of (e) the galaxy-shaping power of black holes began to emerge (eb) with the ability to see plasma.** |
| (1851) **Data from a now-defunct X-ray satellite is providing (eb) new insights into the complex tug-of-war between galaxies, the hot plasma that surrounds them, and the giant black holes that lurk in their centres.** |

(1852) **Launched from Japan on February 17, 2016, the Japanese space agency (JAXA) Hitomi X-ray Observatory functioned for just over a month before contact was lost and (e&eb) the craft disintegrated.**

(1853) **But the data obtained during those few weeks was enough to paint (eb) a startling new picture of the dynamic forces at work within galaxies.**

(1854) **New research, published in the journal Nature today, reveals data that shows just how important the giant black holes in galactic centres are to the evolution (eb) of the galaxies as a whole.**

(1855) **"We think that supermassive black holes act(e&eb) like thermostats," said Brian McNamara, University Research Chair in Astrophysics at the University of Waterloo. "They regulate (e&eb) the growth of galaxies."**

**Champagne bubbles of plasma**

(1856) **During its brief life, the Hitomi satellite collected X-ray data from (e) the core of the Perseus cluster, an enormous gravitationally-bound grouping of (e) hundreds of galaxies.**

(1857) **Located some 240 million light years from earth, the Perseus cluster is one of (e) the largest known structures in the universe.**

(1858) **The cluster includes not only the ordinary matter that makes up the galaxies, but an (e) "atmosphere" of hot plasma with a temperature of tens of millions of degrees, as well as a halo of (e) invisible dark matter.**

(1859) **Earlier studies, going back to the 1960s, have shown that each of the galaxies in the cluster - and indeed most galaxies - likely contains (e) a supermassive black hole in its centre, an object 100 million to more than ten billion times as massive as our sun**

**New look at the galaxy-shaping power of black holes**

(1860) **Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving.**

(1861) **"These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science.**

(1862) **"Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy."**

(1863) **This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return.**

(1864) **The remaining matter gets absorbed into the black hole, adding to (e) its mass.**

(1865) **The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne.**

(1866) **The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution.**

(1867) **Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The**

**answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form.**

(1868) **"Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara.**

(1869) **"It creates (eb) these bubbles, and the bubbles keep the plasma hot.**

(1870) **That's what prevents galaxies from (e) becoming even bigger than they are now."**

(1871) **Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.**

(1872) **In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.**

(1873) **Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.**

(1874) **"It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.**

(1875) **"It's like nature's thermostat, which keeps these galaxies from (e) growing.**

(1876) **If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."**

**A new look at the galaxy-shaping power of black holes**

(1877) **An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the**

(1878) **McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.**

(1879) **Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.**

**Doomed satellite missions**

(1880) **Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.**

(1881) **At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.**

(1882) **Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a**

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| **leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.** |
| (1883) **The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.** |
| (1884) **These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.** |
| (1885) **"But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people** |
| (1886) **. It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."** |
| (1887) **The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.** |
| (1888) **"This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website** |

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

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| (1889) **Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.** |
| (1890) **They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.** |
| (1891) **However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.** |
| (1892) **Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.** |
| (1893) **X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well.** |
| (1894) **The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma.** |
| (1895) **These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6.** |
| (1896) **Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus.** |

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| (1897) | **A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate.** |
| (1898) | **Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure.** |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (1899) | **This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy.** |
| (1900) | **Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole.** |
| (1901) | **But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy.** |
| (1902) | **A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole.** |
| (1903) | **It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth.** |
| (1904) | **"There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA).** |
| (1905) | **This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies.** |
| (1906) | **However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated.** |
| (1907) | **In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1.** |
| (1908) | **We know dark matter exists only from (e) its gravitational effects.** |
| (1909) | **It holds together galaxies and (e&eb) galaxy clusters.** |
| (1910) | **Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years.** |
| (1911) | **To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies.** |
| (1912) | **They used star motions as tracers to (e) weigh the galaxies' central black holes.** |
| (1913) | **X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto.** |
| (1914) | **They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone.** |
| (1915) | **This connection is likely to be related to (e&eb) how elliptical galaxies grow.** |

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| (1916) | **An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together.** |
| (1917) | **Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole.** |
| (1918) | **"In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp** |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (1919) | **The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the …** |
| (1920) | **Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness?** |
| (1921) | **Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness.** |
| (1922) | **"Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain."** |
| (1923) | **Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness.** |
| (1924) | **"With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication."** |
| (1925) | **Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain** |

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| **changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories.** |
| (1926) **The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen.** |
| (1927) **In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer.** |
| (1928) **Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants.** |
| (1929) **These were treated as (=) "aware" and "unaware" trials, respectively.** |
| (1930) **Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it.** |
| (1931) **The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another.** |
| (1932) **Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion.** |
| (1933) **No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness.** |
| (1934) **"We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said.** |
| (1935) **"Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness."** |
| (1936) **The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner.** |
| (1937) **Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain.** |
| (1938) **"We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said.** |
| (1939) **"This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients** |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

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| (1940) | **Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter.** |
| (1941) | **As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined.** |
| (1942) | **Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions.** |
| (1943) | **"Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.** |
| (1944) | **The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.** |
| (1945) | **However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.** |
| (1946) | **The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.** |
| (1947) | **Noda and his colleagues instead studied the change in emissions over (e&eb) time.** |
| (1948) | **By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.** |
| (1949) | **Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).** |
| (1950) | **The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.** |
| (1951) | **Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.** |
| (1952) | **"The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp** |

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(1953) **The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking "palm tree" organization.**

(1954) **The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society**

(1955) **(Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible?**

(1956) **With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem.**

(1957) **This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other.**

(1958) **According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible.**

(1959) **This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges.**

(1960) **In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century.**

(1961) **Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching.**

(1962) **In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom.**

(1963) **Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth.**

**Spheres on trees**

(1964) **Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration.**

(1965) **They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement.**

(1966) **Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum.**

(1967) **By visualizing the problem in this way, the researchers noticed that these particular graphs don't (e) have lots of branches extending from lots of (e) other branches.**

(1968) **Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure.**

(1969) **The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network.**

(1970) **As a result, it provides (eb) an important clue for finding the global minimum.**

(1971) **It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum.**

(1972) **This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large.**

**Thompson's problem 2**

(1973) **By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society**

(1974) **The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima.**

(1975) **The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node.**

**Implications for Thomson's and Smale's problems**

(1976) **Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before.**

(1977) **"Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around."**

(1978) **As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem.**

(1979) **"I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view."**

(1980) **In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges.**

(1981) **From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear.**

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| (1982) | **"While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK.** |
| (1983) | **"For larger numbers, there will likely be (=) multiple funnels.** |
| (1984) | **Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum."** |
| (1985) | **Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.** |
| (1986) | **"The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.** |
| (1987) | **"Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org** |

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**Introduction:**

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price, University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss.**

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| **(289)** | **Black Holes are regions of (e) space in which gravitational fields are so strong that no particle or signal can escape (eb) the pull of gravity.** |
| **(290)** | **The boundary of this no-escape region is (=) called the event horizon, since distant observers outside the black hole cannot (e) see (cannot get light from) events inside.** |
| **(291)** | **Although the fundamental possibility of such an object exists within (eb) Newton's classical theory of gravitation, Einstein's theory of gravity makes (eb)black holes inevitable under some circumstances.** |
| **(292)** | **Prior to the early 1960s, black holes seemed to be (=) only an interesting theoretical concept with no (e) astrophysical plausibility, but with (e&eb) the discovery of quasars in 1963 it became clear that very exotic astrophysical objects could exist.** |
| **(293)** | **Nowadays it is taken for granted that black holes do exist in (eb) at least two different forms.** |
| **(294)** | **Stellar mass black holes are (=) the endpoint of the death of some stars, and supermassive black holes are the result of (e) coalescences in the centers of most galaxies, including our own.** |
| **(295)** | **No signal can propagate from (e) inside a black hole, but the gravitational influence of a black hole is (=) always present.** |

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| **(296)** | **This influence does not (e) propagate out of the hole; it is permanently present outside, and depends only on (=) the total amount of mass, angular momentum, and electric charge that have gone into forming (eb) the hole.** |
| **(297)** | **Black holes can be detected through the influence of this (e&eb) strong gravity on the surroundings just outside the hole.** |
| **(298)** | **In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them.** |
| **(299)** | **Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors.** |

**Classical vs. relativistic black holes**

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| **(300)** | **Something like a black hole exists within (eb) Newton's classical theory of gravity.** |
| **(301)** | **In that theory, an energy argument tells us that there is (=) an escape velocity $v_{esc}=\sqrt{2GM/R}$ from (e) the surface of any spherical object of (e) mass M and radius R** |
| **(302)** | **If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity.** |
| **(303)** | **Thus the condition for such an "unseeable" object is $R<2GM/c2$** |
| **(304)** | **In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape.** |
| **(305)** | **This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity.** |
| **(306)** | **To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it.** |
| **(307)** | **The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c .** |
| **(308)** | **This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed.** |
| **(309)** | **An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity.** |
| **(310)** | **Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable.** |

**Stationary black holes**

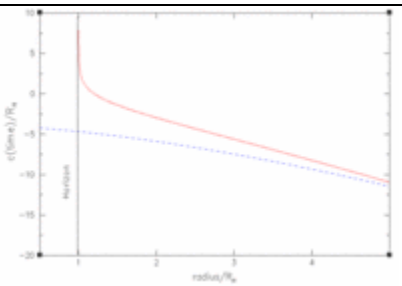
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| **(311)** | **In the Newtonian theory, gravity is described by (e) the potential $\Phi$.** |
| **(312)** | **Inside a spherical object the form of (e) $\Phi(r)$ depends on (=) the interior structure, but in the vacuum outside (e) matter the potential $-GM/r$ depends only on (=) the interior mass.** |
| **(313)** | **Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object.** |

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| **(314) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward.** |
| **(315) This horizon radius RH=2GM/c2 is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory.** |
| **(316) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution.** |
| **(317) Radius RHhere actually means that the area of the event horizon is (=) 4πR2H.** |
| **(318) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution.** |
| **(319) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity.** |
| **(320) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own.** |
| **(321) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole.** |
| **(322) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density.** |
| **(323) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed.** |
| **(324) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole.** |
| **(325) In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter.** |
| **(326) Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum.** |
| **(327) In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior.** |
| **(328) The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object.** |
| **(329) Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum.** |
| **(330) These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.]** |
| **(331) All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time.** |
| **(332) But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?"** |
| **(333) The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point.** |

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| **(334) One of these coordinates is (=) called the "coordinate time."** |
| **(335) Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate.** |



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| **(336) Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole** |
| **(337) At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole.** |
| **(338) Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations.** |
| **(339) For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree** |
| **(340) An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events.** |
| **(341) The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving.** |
| **(342) Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle.** |
| **(343) The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return.** |
| **(344) But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon.** |
| **(345) Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future.** |
| **(346) The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?"** |

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| **Black hole parameters** |

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| **(347) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin).** |
| **(348) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun.** |
| **(349) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits.** |
| **(350) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form.** |
| **(351) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds.** |
| **(352) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2.** |
| **(353) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed.** |
| **(354) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer.** |
| **(355) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2)** |
| **(356) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials.** |
| **(357) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable** |
| **(358) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite.** |
| **(359) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole.** |
| **(360) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars.** |
| **(361) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity.** |
| **(362) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity.** |
| **(363) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite.** |
| **(364) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity.** |
| **(365) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum.** |

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| **(366) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship.** |
| **(367) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= (4πG2/c4)[(M+M2−a2−−−−−−−√)2+a2].** |
| **(368) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored.** |

**Dynamical black holes**

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| **(369) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape.** |
| **(370) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, a horizon can be born.** |
| **(371) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid.** |
| **(372) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward.** |
| **(373) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations.** |
| **(374) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes.** |
| **(375) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3.** |
| **(376) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/** |
| **(377) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers.** |
| **(378) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission.** |
| **(379) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.** |
| **(380) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.** |
| **(381) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.** |
| **(382) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.** |

**(383) The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature 6×10−8 (Mo/M) K.**

**(384) Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.**

Then $[Y_6]_6[\gamma_6]_6, a_6) = E_6(a_6)e^{-\gamma_6 a_6}$. by a simple calculation,

$$\int_0^\infty \beta_6(a_6)Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma}\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}\ da\right);$$

$$\int_0^\infty Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma_6}\left(\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}\ da_6\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_6(0)}{\gamma_6}[f_{60}^6]\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}da_6\ (\int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6)$$

$$+\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right).$$

$$\left(1 + \frac{(\rho_0)_6(0)}{\gamma_6}(f_{60}^{'})_6\left(\int_0^\infty E_6(\,a_6)da_6 - \int_0^\infty E_6(\,a_6)e^{-\gamma_6 a_6}da_6\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right)(1 + \frac{(f_{60}^{'})_6}{\gamma_6}(\rho_0)_6(0)\int_0^\infty E_6(\,a_6)da_6) = 0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1 - \int_0^\infty \beta_6(a_6) Y_6(\gamma_6 a_6) da_6 = 1 - \int_0^\infty \beta_6(a_6) E_6(a_6) e^{-\gamma_6 a_6} da_6 \neq 0,$$

Since $\int_0^\infty \beta_6(a_6) Y_6(\gamma_6 a_6) \int_0^a \alpha_6(\tau_6) Y^{-6}(\gamma_6, \tau_6) d\tau_6 da_6 > 0$. Therefore, the characteristic solution is

$$\gamma_6 = -(f_{60}^{'})_6 (\rho_0)_6 \int_0^\infty E_6(a_6) da_6.$$

Since $f_6'(P_0)_6 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_6), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE OF DEBIT TYPES AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_6), 0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_6 + (\mu_6(a_6) + [f_6]_6(P_6^*) + [f_2]_6(Q_6^*))P_6^* = 0, \quad (26)$$

$$p_6^*(0) = h_6(Q_6^*) \int_0^\infty \beta_6(a_6) p_6^*(a_J) p_6^*(a_6) da_6 \quad (27)$$

$g_6(P_6^*) - b_6 = 0,$ (28)

Where $P_6^*$=$\int_0^\infty p_6^*(a_6)\, da_6$

Equation (28) yields $P^* = g^{-6}(b)$,

From (26),

$p_6^*(a) = p_6^*(0)e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6}$

Substituting into (27), we need to solve

$$\left(1 - h_6(Q_6^*)\int_0^\infty \beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6} da_6\right) p_6^*(0) = 0$$

For $Q_6^*$

Fix $p_6^*(0)$ and $P_6^*$ Let

$$H_6(z_6) = 1 - h_6(z_6)\int_0^\infty \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6} da_{,6}$$

And assume that

$$[k_3]_6 \int_0^\infty \beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(g_6^{-6}(b_6))a_6 - [k_2]_6 a_6}\, da_6 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_6(z_6)] > 0.$$

Clearly,

$$\frac{dH_6}{dz_6} = h_6(z)([f_2^{'}]_6(z_6))\int_0^{\infty} a_6\beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - f_6(P_6^*)a_6 - f_2(z_6)a_6}\, da_6$$

$-h^{'}(z_6)\int_0^{\infty}\beta_6(a_6)\, e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6}\quad da_6$

By the assumptions on $\left[f_2^6\right]_6 and\ [h^6]^6 dH_6/dz > 0$. Hence, if

$$H_6(0) = 1 - \int_0^{\infty}\beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6}da_6 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_6(0) < 0$ follows if $g_6^{-}(b_6) < [f_6^{-6}]_6(r_6)$, where $r_6$ is the solution of (12).

Indeed, since

$$R_6(s_6) = \int_0^{\infty}\beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 e^{-y}6 a}da_6$$

Is a decreasing function of $s_6$ when $P_6^* = [g^{-6}]_6(b_6) < [f_6^{-6}]_6(r_6)$, that is $[f_6]_6(P_6^*) < r$, $R\left(f_6(P^*)\right) > 1$which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_6(g^{-6}(b))) > 6, i.e.,$

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6\left(g_6^{-6}(b)\right)a_6} > 1. \quad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_6^*(a_6), Q_6^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_6^*(a_6), Q_6^*)$ is more complicated. We again consider perturbations $m_6(t_6 a_6) = p_6(t_6, a_6) - p_6^*(a_6)$ and $W_6(t) = Q_6(t) - Q^*$.Then

$$M_6(t) \equiv \int_0^\infty m_6(t, a_6) da_6 = P_6(t) - P_6^*,$$

$$[m_6]_6 + [m_a]_6 = [p_6]_6 + [p_a]_6 - \left[p^*\right]_6(a_6) \approx -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6$$

$$-[f_{6*}^6]_6 p_6^*(a_6)M - [f_{2*}^6]_6 p_6^*(\mathrm{a})W_6,$$

$$m_6(t,0)p_6(t,0) - p_6^*(0) \approx h_* \int_0^\infty \beta_6(a_6) m_6(t, a_6)\, da + \frac{[h_*^6]_6}{[h_*]_6}\, p_6^*(0) W_6$$

$$\frac{dW_6}{dt} = \frac{dQ_6}{dt} \approx [g_*^6]_6 Q_6^* M_{,6}$$

Where $[f_{6*}]_6$ denoting $[f_6(P^*)]_6$, $[f_{6*}]_6$ denoting $[f_2(P^*)]_6$, $\left[f_{6*}^{'}\right]_6$ denoting $\left[f^{'}{}_6(P^*)\right]_6 \left[f_{2*}^{'}\right]_6$ denoting $[f_2]_6$ $(Q_6^*,$ $h_{*}^{}$ denoting $h_*^{'}$ denoting $h_6^{'}$ $(Q_6^*)$, and $\left[g_*^{'}\right]_6$

Denoting $\left[g^{'}\right]_6(P_6^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_6]_6 + [m_a]_6 = -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6 - [f'_{6*}]_6[p_6^*](a_6)M_6 - [f'_{2*}]_6 p_6^*(a_6)W_6$$

$$m_6(t,0) = [h_*]_6 \int_0^\infty \beta_6(a_6)m_6(t,a_6)da_6 + \frac{[h_*^6]_6}{[h^*]_6}[p^*]_6(0)\,W_{,6} \qquad (30)$$

$$\frac{dW_6}{dt} = [g_*^6]_6[Q_6^*]M_6$$

Integrating $[m_6]_6(t,a_6)$ with respect to a from 0 to ∞, we have

$$\frac{dM_6(t)}{dt} = \int_0^\infty [mt]_6(t,a_6)da_6$$

$$= -\int_0^\infty [m_a]_6\,(t,a_6)da_6 - \int_0^\infty \mu(a_6)m(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + [f'_{6*}]_6 \int_0^\infty p_6^*(a_6)da_6\right)M_6 - [f_{2*}^6]_6 \int_0^\infty p_6^*(a_6)da_6 W_6$$

$$= m_6(t,0) - \int_0^\infty \mu_6(a_6)\, m_6(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + [f_{6*}']_6 P_6^*\right)M_6 - [f_{2*}']_6 P_6^* W_{,6}$$

From (30), we arrive at the equations

$$\frac{dW_6(t)}{dt} = \int_0^\infty ([h_*]_6\beta_6(a_6) - \mu_6(a_6))m_6(t,a_6)da_6$$

$$-\left([f_{6*}] + [f_{2*}]_6 + [f_{6*}']_6 P_6^*\right)M_6 + \left(\frac{[h_*']_6}{[h_*]^6}\,P_6^*(0) - f_{2*}' P_6^*\right)W_6,$$

(36)

$$\frac{dW_6(t)}{dt}\left[g^{'}{}_{*}\right]_6 Q_6^{*} M_{.6}$$

By the mean value theorem,

$$\int_0^{\infty}([h_{*}]_6\beta_6(a_6)-\mu_6(a_6))m_6(t,a_6)da_6=([h_{*}]_6\beta_6(\theta_6)-\mu_6(\theta_6))M_{,6}$$

Where $0<\ \theta_6=\theta_6(t)<\infty$.

And if $a_6^2\,[a_6^2]_6<4\,[a_2]_6<4\,a_2 e$,

$\mathrm{Re}\,\lambda_6\left[\lambda_6\right]_6=\mathrm{Re}\lambda_2\left[\lambda_2\right]_6=-\frac{6}{2}a_6[a2\ ]_6<0.$

Let $\Delta=\max_{2<u<\infty}\left|h_{*}\,[h_{*}]_6\beta_6(a_6)-\mu_6(a_6)\right|$. Then$|n_6(t)|$ 2Δ∀t 0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^{*}(a),\boldsymbol{Q}^{*})$ is asymptotically stable if either $a_6^2<4\,a_2 e$ and $[a_6]>2\Delta$,or $[a_6^2]$ 4$[a_2]$e and $(a_6-\sqrt{(a_6^2)_6-4(a_2)_6 e})>2\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a)=\mu_0+\left(\frac{a}{a+6}\right)$ and $\beta_6(a_6)=ae^{-ma}$ Then $\Delta=[\mu_0]+1$ Simple calculation leads to

$$\int_0^{\infty}\beta_6(a_6)e^{-\int_0^a\mu_6(\tau_6)+d\tau_6}da=\frac{6}{(m+\mu_0+1)^2}+\frac{2}{(m+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_6 + [\mu_0]$ 0.5. Assume$[f_6](x_6)= K_6x_6$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_6(x_6)= [0x_6]/(1 + x_6)$; the functional responses $[[f_2](x) = [0]\, x_6/(1 + x_6)$and $h_6(x_6)$=(1+$\omega_6 x_6$)/$(1 + x_6)$(See Holling [7] and Rashevsky [13])

Equation (64),now has the form of

$$\frac{m_6 + [\mu_0] + [\mathrm{K}][P_0])}{[m_6]_{\,6} + (\mu_0)_{\,6} + [\mathrm{K}][P_0])\,{}^{3}_{6}\ [\mathrm{K}]_{\,6}[P_0]\,{}^{3}_{6}} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(6 + P_0) < b_6$i e $[P_0] < b/(\delta_6 - b_6)$then the equilibrium$(P_0)$ $[P_0]_{\,6}(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_6$

$[f_{6*}] = K_6\ \mathrm{P}^{*}\ [f_{2*}] = \frac{\sigma \mathrm{Q}_6^*}{6+\mathrm{Q}_6^*}$ $[h_*] = \frac{6+\omega_6 \mathrm{Q}_6^*}{6+\mathrm{Q}_6^*}$

$[f_{6*}] = K_6\ \left[\mathrm{f}\,{}^{6}_{2*}\right] = \frac{\sigma_6}{(6+\mathrm{Q}_6^*)^2}[\ \mathrm{g}\,{}^{6}_{*}] = \frac{\sigma_6}{(6+\mathrm{P}_6^*)^2}$ $\left[\mathrm{h}\,{}^{6}_{*}\right] = \frac{\omega_6 - 6}{(6+\mathrm{Q}_6^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0,$ K, b, δ, and $\omega$ need to satisfy

$$\frac{\omega_6}{(m + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 + b_6} + \sigma)^{\,2}} + \frac{2\omega_6}{([m_6]_{\,6} + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 + b_6} + \sigma_6)^{\,3}} < 1$$

And

$$\frac{1}{(m_6 + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 - b_6})^{\,2}} + \frac{2\omega_6}{(m_6 + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 - b_6})^{\,3}} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_6 + [\mu_0]_6 + 1 + \frac{K_6 b_6}{\delta_6 - b_6} + (m_6 + [\mu_0]_6 + \frac{K_6 b_6}{\delta_6 - b_6} + \sigma_6) Q_6^*)^3$$

$$= m_6 + [\mu_0]_6 + 3 + \frac{K_6 b_6}{\delta_6 - b_6}(1 + Q^*)(1 + \omega_6 Q_6^*$$

Now

$$\left[a_6^2\right] - 4\left[a_2\right]e = \left(\frac{2K_6 b_6}{\delta_6 - b_6}\right)^2 + \frac{4K_6 b_6 \sigma_6 Q_6^*}{(\delta_6 - b_6)(1 + Q_6^*)} + \frac{\sigma_6^2 Q_6^{*2}}{(1 + Q_6^*)^2}$$

$$\frac{4 b_6 \sigma_6 (\delta_6 - b) Q^*}{\delta_6 (1 + Q_6^*)^2} - \frac{4 (\delta_6 - b_6)^2 (1 - \omega_6) Q_6^*}{\delta_6 (1 + \omega_6 Q_6^*) 1 + Q_6^*)} P * (0)$$

And

$$\Delta = \lim_{0 < a < \infty} |[h_*]_6 a e_6^{mu} [\mu_0]_6 - \frac{a_6}{a_6 + 6}$$

Hence if $[a_6^2 > 4(a_2)_6 e$ and

$$[a_6] = \frac{2K_6 b_6}{\delta_6 - b_6} + \frac{\sigma_6 Q_6^*}{6 + Q_6^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_6]_6 = [f_{6*}]_6 + [f_{2*}]_6 + [f_{6*}]\,P^*_6$$

$$[a_2]_6 = \left[f^6_{2*}\right]P^*_6 - \frac{h^6_*}{h_*}P^*(0)$$

$$e=[g^6_*]_6\,Q^*_6$$

$$\eta_6(t) = \left[h_*\right]_6\beta_6\left(\theta_6(t)\right) - \mu_6\left(\theta(t)\right),$$

The system(36)can be written as the following system:

$$[M(t)]\left[a_6\right]M_6 - [a_2]_6 W_6 - [n(t)]M_6$$

$$w_6(t) = eM_6$$

Write

$$A=\begin{pmatrix} -[a_6]_6 - [a_2]_6 & \\ e & 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_6 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_6 \\ W_6 \end{pmatrix}$$

The solution of (32)is

$$X_6(t) = e^{A,t}\,[X]_6(0) + \int_0 e^{At}$$

Other problems to which the model is applied are: To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case β→∞ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness condition of (e) entropy. It is shown that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology

(gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure authors have discovered. It is shown that Legendre invariant thermodynamic scalar curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws, leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions,**

**geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged AdS black holes by considering the case of the RN-AdS black holes surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase transition. For the case $\omega_q=-2/3$ authors derive analytic expressions of (e) critical physical quantities, while (e) for cases $\omega_q\neq-2/3$ appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark

energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=)

chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross

section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions, while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center

collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma? Paul M. Chesler STOP Paul M. Chesler, Krishna Rajagopal** calculate how the energy and (e&eb) the opening angle of jets in N=4 SYM theory evolve as they propagate through (e&eb) the strongly coupled plasma of that theory. **Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb)

heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version) **Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb) strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion

collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar (=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI: 10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills

(SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the

decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed

black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to

equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a

nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about

the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic

map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstates thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11 August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical

advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.**

Characterizing the behaviour of strongly coupled quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at $\lambda=\lambda c$, the correlation length diverges as $\xi\sim|\lambda-\lambda c|-\nu$, while (e) the gap between the ground state and first excited state

vanishes as (e) ΔE∼ξ−z∼|λ−λc|zν. Exactly at the critical point λ=λc, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as ω∼|k−k0|z. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004 STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting

qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element. The basic architecture presented here can be expanded, offering (eb) the possibility for the coherent interaction of (e&eb) a large number of superconducting qubits. **Nature 449, 438-442 (27 September 2007) | doi:10.1038/nature06124; Received 18 April 2007; Accepted 25 July 2007 Coherent quantum state storage and transfer between two phase qubits via a resonant cavity Mika A. Sillanpää1, Jae I. Park1 & Raymond W. Simmonds1 STOP. Chui-Ping Yang, Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e)

long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007**

# PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION ATTRITION MODELS

## MOODULE NUMBER SEVEN

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Systems to which have applied the model are as follows: Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures have been proposed. An attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in**

**(eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP. In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011 http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 Qubit transfer**

between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys. Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer. Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP. In recent years S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P have investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys.

**Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP. Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2 [quant-ph] for this version) Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP. Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue from over (e&eb) a thousand plant and animal samples. NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP. Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this behavior, enhancing and accelerating (eb, eb+) the growth of correlations. Michael Kastner gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. Michael Kastner tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP. A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E.**

**Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91, 075435 – Published 26 February 2015 STOP. Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1 proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point'' with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. Membranes at quantum criticality Petr Hořava1, 2 Published 4**

**March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that (eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035 Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013 STOP Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study**

**the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen STOP Authors study the dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing**

**seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) Dynamic Service Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by H∗(X,(Λ∗TX)⊗g), where X is the target hyperk\"ahler manifold. We further give a mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the**

**low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition. Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb) the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537 © 2016 American Physical Society Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant (Δα/α) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and Δα/α to (e&eb) this conclusion.**

**However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts z≲3, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the Λ cold dark matter model value (jΛCDM=1); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in α(z), linked to (e&eb) a linear coupling (ζ) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order |ζ|∼10−4, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in Δα/α (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems, authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them**

**are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP Authors measure the star formation properties of (e) two large samples of galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the star formation rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a function of environment. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]. Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size √ n × √ n? The problem is that √n time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time O (√n) for d ≥ 3, or O (√nlog5/2n) for d = 2. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ∗This work was mostly done while the author was a PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship and by CIAR. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In**

**d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e) sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP The relationship between the foundations of mathematics and (e&eb) physics is a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb)**

**flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph] (or arXiv: 1503.05600v1 [math-ph] for this version) Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative**

**real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to the fine structure constant. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor goes to 0 as the (e&eb) time goes to 0, the big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph]**

**(or arXiv: 1306.4613v1 [quant-ph] for this version) Effects of gauge theory based number scaling on geometry Paul Benioff STOP In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x, outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph] (or arXiv: 1211.3381v2 [quant-ph] for this version) Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value, \int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances, the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an**

**observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include (e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP This work is based on (e) the field of reference frames based on (e) quantum representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul Benioff STOP A quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of**

**these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference: Journal of Physics: Conference Series 70 Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows: System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The**

**Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol**

## 32. MODEL FORMULATION

Let $p_7(t, a_7)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(g_7)Q_7$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_7(0) = 0, \quad g_7(P_7) > 0, \ P_7 \ 0, \qquad \lim_{P_7 \to \infty} g_7(P_7) = (K_7)_7 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_7 + (\rho_a)_7 + D_7(a_7, P_7, Q_7)\rho_7 + 0, \qquad 0 < a_7 \ \infty,$$

$$\rho_7(t,0) = \int_0^\infty n_7(a_7, P_7, Q_7)\rho_7(t_7, a_7)da_7$$

Where $D_7$the mortality or debit type function is function and $n_7$ is the fertility function.

Put $_7$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_7}{dt}\right]_7 = x_7\left[\frac{dQ_7}{dt}\right]_7 - x_7\left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_7}{dt}\right]_7 = x_3\left[\frac{dQ_3}{dt}\right]_7 - x_4\left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_7}{dt}\right]_7 - x_6\left[\frac{dQ_3}{dt}\right]_7$$

Where $x_7, x_7, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_7$ and $g_7(P_7)Q_7$

Now when you add the term $(b_7Q_7)$ with $g_7(P_7)(Q_7)$ we are infect see is that $(b_7Q_7)$ is credited incrementally part of $g_7(P_7)(Q_7)$ or instant nearly, when subtracted from $\frac{dQ_7}{dt}$ , it means that $g_7(P_7)(Q_7)$ is slowly eaten by $\frac{dQ_7}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

Following equation throws light on important issues in the following cases

**Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA**

(1988) (Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor.

(1989) In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface.

(1990) As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone.

(1991) Such planets are all the "right" distance from (e) their star, and hopefully, also have water.

(1992) But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone.

(1993) For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement.

(1994) But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway.

(1995) But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain.

(1996) It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life.

(1997) Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface.

(1998) The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113

**Abstract**

(1999) The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question.

(2000) Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry.

(2001) This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry.

(2002) Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments.

(2003) Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy.

(2004) The electronic and (e&eb) structural variability among them is extraordinary.

(2005) The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations.

(2006) The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface.

(2007) The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers.

(2008) Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed

mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp

**Water waves travel back in time to retrace their ripples By Jacob Aron (New Scientist)**

(2009) Nudge your bath tub just right and (eb) you might end up reversing time.

(2010) Don't worry – you won't end up wet, naked and without a towel in the Cretaceous period. Emmanuel Fort at the Langevin Institute in Paris, France, and his colleagues have developed what they call a "time mirror", which can make water waves rewind into (e&eb) their past. Humans and rubber duckies are (=) affected.

(2011) To understand how this works, imagine dropping a stone in a pond and watching waves ripple out from a central point in an ever-expanding circle. The equations governing the behaviour of waves say that it should also be possible for a ring of ripples to refocus back into a single point, like a reversed movie of the dropped stone. But we never actually see this happening, because it is vanishingly unlikely that an outside-in ring would form spontaneously. Put it in reverse

(2012) Fort and others have previously shown that it is possible to engineer the situation so the waves turn back on themselves by kitting out the edge of your pond with sensors and wave generators, effectively recording and then playing back the waves in reverse. The process has been done with water, sound and even microwaves, and is used to produce clearer ultrasound images for medical uses.

(2013) But now the team has discovered that simply dropping the tub downwards is enough to send water waves scurrying in reverse, with no need for complex sensors. That's because the velocity of water waves depends on acceleration due to gravity. "If you do a sudden jolt, you change the effective gravity, and the velocity of the wave changes," says Fort.

(2014) The disruption essentially resets the starting point of the wave. Once normal gravity is restored, it starts to travel both backwards and forwards. "Suddenly, you freeze the wave. When you release it again, it doesn't remember that it was going one way," says Fort.

(2015) Ripples in the water returned exactly to their start point - the shape of the Eiffel Tower

(2016) Ripples in the water returned exactly to their start point – the shape of the Eiffel Tower

**Déjà vu**

(2017) To show that the backward-travelling waves converge exactly on their starting point, just as if you were rewinding time, the team conducted a series of experiments. In some, they simply created a circular ripple of waves by dipping a probe in the water, similar to dropping a stone in the pond. In others, they used wave emitters in the shape of France, the Eiffel Tower and a smiley face to create complex wave-fronts. In each case, the reversed waves converged to recreate their original shape.

(2018) The experiments aren't truly reversing time, because the waves continue to travel out again once they have converged at the central point, points out Fort. "A real time-reversal would have a sink that absorbs all the waves," he says.

(2019) Now the team hopes to conduct similar experiments with other kinds of waves, like light or sound, although it may prove more difficult to find an analogue of their time-mirror jolt in other wave mediums, he says.

(2020) "The work is spectacular," says Geoffrey Edelmann of the US Naval Research Laboratory in Washington, D.C. "They essentially change the rules that govern how the wave propagates and, in doing so, create an entirely new way to focus waves free of spatial constraints." He also thinks the technique could be extended to other kinds of waves, and perhaps have medical applications such as lithotripsy, which uses ultrasound waves to break up kidney stones. Journal reference: Nature Physics, DOI: 10.1038/nphys3810

**Time reversal and holography with spacetime transformations Vincent Bacot, Matthieu Labousse, Antonin Eddi, Mathias Fink & Emmanuel Fort Nature Physics (2016) doi: 10.1038/nphys3810**

(2021) Wave control is usually performed by (e) spatially engineering the properties of a medium.

(2022) Because time and space play similar roles (e&eb) in wave propagation, manipulating time boundaries provides (eb) a complementary approach.

(2023) Here, authors experimentally demonstrate the relevance of this concept by (e) introducing instantaneous time mirrors.

(2024) Authors show with water waves that a sudden change of (e) the effective gravity generates (eb) time-reversed waves that refocus at the source.

(2025) They generalize this concept for (e) all kinds of waves, introducing (eb) a universal framework which explains (eb) the effect of any time disruption on (e&eb) wave propagation.

(2026) Authors show that sudden changes of the medium properties generate (eb) instant wave sources that emerge (eb) instantaneously from the entire space at (eb) the time disruption.

(2027) The time-reversed waves originate from (e) these 'Cauchy sources', which are (+) the counterpart of Huygens virtual sources on (eb) a time boundary. It allows us to revisit the holographic method and introduce a new approach for wave control.

**Viewpoint: A Bird's Eye View of Circuit Photons Sebastian Schmidt, Institute for Theoretical Physics, ETH Zurich, 8093 Zurich, Switzerland June 30, 2016• Physics 9, 74 A scanning probe detects the quantum states of photons in a microwave circuit, providing the information needed for quantum simulations. Figure caption expand figure A. Houck/Princeton (deleted)**

(2028) Figure 1: Scanning defect microscopy provides a map of photons in a resonator lattice. Houck and colleagues demonstrated the technique using (e) 49 resonators (grey lines) that were coupled together to (e&eb) form a kagome lattice.

(2029) Quantum mechanics rules the dynamics of (e) light and matter.

(2030) Yet performing a quantum-mechanical simulation of a material from first principles is (=) practically impossible on a classical computer because (e) the complexity of the simulation increases (eb) exponentially with the number of particles involved.

(2031) The solution, according to Richard Feynman, was to build a machine out of (e) quantum building blocks that could directly emulate (e&eb) the material itself [1].

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| (2032) Prototypes of such quantum simulators that are based on (e) ultracold atoms, ions, photons, and superconducting microwave circuits are now available [2], with (e&eb) the latter, in particular, having attracted (e) Silicon Valley's interest. |
| (2033) The challenge with these circuit-based simulators, however, is that they are (=) 2D, which complicates the readout of their constituent elements. Andrew Houck from Princeton University, New Jersey, and colleagues have now delivered an attractive solution by developing a technique [3], called scanning defect microscopy (Fig. 1), that determines (eb) the number of photons occupying (e) each mode of a 2D microwave circuit. |
| (2034) It is this information that would serve as (=) the fundamental input and output for (e) certain quantum simulations. |
| (2035) Superconducting microwave circuits combine electronic and (e&eb) photonic degrees of freedom [4, 5]. |
| (2036) The main element of the circuit is (=) a transmission line, which is made up of (e) a central superconducting wire separated by a gap from (e) two grounded plates. |
| (2037) All of these structures are on a single plane, as if one had taken a 2D slice through (e&eb) a coaxial cable. |
| (2038) When truncated, the transmission line becomes (=) a resonator, which can host (e&eb) discrete photon modes within (eb) its gaps |
| (2039) Large lattices of resonators can be engineered in various 1D or 2D geometries by coupling (e&eb) two, three, or more resonators together via (e&eb) a capacitive interface. In many ways, photons in such devices behave similarly to electrons in a solid. |
| (2040) To make microwave circuits that can simulate (e&eb) quantum phenomena faster than a classical computer, however, resonator lattices have to be integrated with (e&eb) superconducting qubits |
| (2041) . Such qubits are controlled with (e&eb) electrical currents in a Josephson tunnel junction, and in many respects, they behave (e&eb) like artificial atoms, which couple to (e&eb) the photons in the circuit. |
| (2042) As a result, superconducting microwave circuits can be used to (e) explore the coupling between the quantum states of light and (e&eb) matter, the regime of (e) circuit quantum electrodynamics (cQED). |
| (2043) Photons in these devices often exhibit (eb) striking matter-like behavior [6, 7], providing (eb) the basis for the simulation of complex materials. |
| (2044) Such circuits can be fabricated on a substrate using (e) standard lithographic techniques, with (e&eb) qubits and resonators that are (=) hundreds of micrometers or even millimeters in size. |
| (2045) Small-scale realizations of cQED lattices have already been used for (e) proof-of-principle quantum simulations [8, 9]. |
| (2046) However, future large-scale simulations will have to be initialized and measured (eb) efficiently—steps known as (=) drive and readout. |
| (2047) This requires connecting the lattices to (e) external circuitry, such as voltage sources, circulators, and amplifiers, which typically reside outside (e) the chip. |
| (2048) So far, this connectivity has been realized through (e&eb) external transmission lines, but these elements can only couple to (e&eb) the resonators at the edges of (e) the lattice because of (e) its 2D nature. |
| (2049) Such a side view provides (eb) only limited information about (e&eb) the quantum state and spatial correlations of photons inside (eb) the bulk resonators, seriously limiting (e) the applicability of circuit QED devices for (e) quantum simulations. |

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| (2050) The probe developed by Houck and colleagues, scanning defect microscopy (SDM), solves (e) this problem by (e) moving outside the plane of the chip to provide (eb) a fuller image of photons within (eb) the chip. |
| (2051) In SDM, a small sapphire tip is positioned a few tens of (e) micrometers above the surface of one of the resonators (Fig. 1). |
| (2052) The presence of this dielectric modifies (e&eb) the evanescent electric field from the circuit photons, which extends (eb) just outside the plane of the chip. |
| (2053) Specifically, the tip increases (eb) the local capacitance of the transmission line, downshifting (e) the local resonator frequency. |
| (2054) This in turn causes (eb) a drop in the transmission of photons through (e&eb) the lattice when (e) it is driven at its nominal resonance frequency. |
| (2055) Houck's team showed that if the sapphire tip has only a weak (linear) effect on the resonator frequency, then (eb) a simple analytic relationship exists between the induced frequency shift, (e&eb) the change in transmission, and (e&eb) the local photonic occupation of each resonator. |
| (2056) This relationship provides (eb) a recipe for detecting the local photon occupancy inside (eb) the lattice. |
| (2057) First, calibrate the probe by determining (eb) the frequency shift for (e) a single isolated resonator as (=) a function of the tip's position. |
| (2058) Second, measure the transmission through (e&eb) the whole lattice with and without (e&eb) the probe. |
| (2059) Finally, use the measured transmission drop with (e&eb) the probe present to find (eb) the photonic occupation of (e) an individual resonator in (eb) the lattice. |
| (2060) A map of the photon-mode occupancies can be obtained by (e) scanning the probe across the whole lattice. |
| (2061) Houck's team used (e) SDM to image the occupancies of (e) the photon modes of (e) 49 coupled resonators arranged in (eb) a special geometry called a kagome lattice (see details in Fig. 1). |
| (2062) (To simplify the measurement, no superconducting qubits were (=) incorporated with (e&eb) the resonators.) |
| (2063) The occupancies they measured were in excellent agreement with (=) theoretical predictions, corresponding to (e&eb) a fidelity of 0.99. |
| (2064) The high fidelity is even more impressive considering the weak signal-to-noise ratio of (e) their technique (which is limited by (e) the intentionally small size of the frequency-shifted region) and (e&eb) the technical challenge of (e) positioning the probe inside (eb) a cryostat for (e) in situ measurements. |
| (2065) While SDM shares similarities with (e&eb) so-called bead-pull experiments for accelerators and scanning gate microscopy for (e) nanostructures, it is truly a (=) new tool for quantum electronics. |
| (2066) First application of SDM is likely to be (=) the simulation of 2D topological phenomena and (e&eb) geometrically frustrated lattices (of which the kagome lattice is (=) one example.) |
| (2067) Both are associated with (e&eb) macroscopically degenerate systems, which are (=) notoriously difficult to simulate on (e&eb) a classical computer. |
| (2068) Such circuits are (=) genuinely open systems (photons can be added or lost). |
| (2069) This combined with (e&eb) their inherent strong light-matter coupling, makes (eb) such quantum simulators ideal candidates for (e) the study of many-body, out-of-equilibrium phenomena, such as (=) thermalization or nonequilibrium phase transitions. |

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| (2070) SDM could also lead to (eb) the discovery of new exotic phenomena, which are (=) unknown in traditional solid-state systems. |
| (2071) In a nutshell, SDM constitutes a new and indispensible tool for (e) future experiments on quantum simulations with microwave photons. This research is published in Physical Review X. |

**Imaging Photon Lattice States by Scanning Defect Microscopy D. L. Underwood, W. E. Shanks, Andy C. Y. Li, Lamia Ateshian, Jens Koch, and A. A. Houck Phys. Rev. X 6, 021044 – Published 30 June 2016**

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| (2072) Microwave photons inside (eb) lattices of coupled resonators and superconducting qubits can exhibit (eb) surprising matter like behavior. |
| (2073) Realizing such open-system quantum simulators presents (eb) an experimental challenge and requires (e) new tools and measurement techniques. |
| (2074) Here, authors introduce scanning defect microscopy as one such tool and illustrate its use in mapping (e&eb) the normal-mode structure of microwave photons inside (eb) a 49-site kagome lattice of coplanar waveguide resonators |
| (2075) Scanning is accomplished by (e) moving a probe equipped with a sapphire tip across (e&eb) the lattice. |
| (2076) This locally perturbs (e&eb) resonator frequencies and induces (eb) shifts of the lattice resonance frequencies, which authors determine by (e) measuring the transmission spectrum. |
| (2077) From the magnitude of mode shifts, author can reconstruct (eb) photon field amplitudes at (eb) each lattice site and thus create (eb) spatial images of the photon-lattice normal modes. Received 4 December 2015 DOI:http://dx.doi.org/10.1103/PhysRevX.6.021044 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society |

| |
|---|
| |
| We assume that |
| $$D_7 = \mu_7(a_7) + (f_7)_7(P_7)_7 + (f_7)_7(Q)_7 \quad (7)$$ |
| $$n_7 = \beta_7(a_7)h_7(Q_7), \quad (3)$$ |
| Where $\mu_7$the age- dependent debit rate whit is $\mu_7(a_7)$ 0 $\forall a > 0$; $(f_7)_7$ is the density- dependent death rate with $f_7(0) = 0$, $(f_7^7)_7(P_7)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_7)_7(P_7) = (k_7)_7$ ∞; $f_7$ is the functional response of creditors to prey with $[f_7]_7(0) = 0, (f_7^7)_7(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_7)_7(Q_7) = [k_7]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^7(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0. |

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_7 + (\rho_a)_7 = -\mu_7(a_7) + (f_7)_7 P_7) + \big(f_7)_7(Q_7)\big)(\rho_7)_7 \quad (4)$$

$$\rho_7(t_7, 0) = h_7(Q_7) \int_0^\infty \beta_7 (a_7)\rho_7(t, a_7)da, \quad (5)$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(P_7)Q_7 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_7$ and $Q_7$ of (4-6) satisfy

$$0 \leq \rho_7(t, a_7) \leq \rho_7^*(t, a_7),$$

$$0 \leq Q_7(t) \leq Q_7^*(t),$$

Where $\rho_7^*(t, a_7)$ and $Q_7^*(t)$, are solutions of

$$(\rho_t^*)_7 + (\rho_a^*)_7 = -(\mu_7(a_7) + (f_7)_7(P_7^*) + (f_7)_7 Q_7^*))\, \rho_7^*, \quad (7)$$

$$\rho_7^*(t, 0) = \int_0^\infty \beta_7(a_7)\, \rho_7^*(t, a_7)da_7 \quad (8)$$

$$\frac{dQ_7^*}{dt} = -b_7\, Q_7^* + g_7(P_7^*)Q_7^*, \quad (9)$$

With $P_7^*(t) = \int_0^\infty \rho_7^*(t, a_7)da_7$

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From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_7^*(t)}{dt} = P_7^*(t)\big(r_7 - (f_7)(P_7^*) - (f_7)_7(Q^*)\big), \quad (10)$$

$$\frac{dQ_7^*(t)}{dt} = Q_7^*(t)(-b_7 + g_7(P_7^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_7(a)e^{-\int_0^a \mu_7(\tau)dt}e^{-ra}\,da_7 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_7^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(t)dt}\,da_7 > 1, \quad (H1)$$

Which guarantees $r_7 > 0$.

*Remark.* If $\beta_7(a) = [\beta_0]_7 a_7 e^{-\alpha_7 a_7}, \mu_7(a_7) \equiv \mu_7$ then the solution of (12) is

$$r_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_7]_7(P_7) = 0, [f_7](Q_7) = \lambda_7 Q_7\ h_7(Q_7) = 1, \quad and\ \ g_7(P_7) = c_7 P_7$$

$$[\rho_7]_7 + [\rho_a]_7 = -(\mu_7 + \lambda_7 Q_7)\rho_7$$

$$\rho_7(t,0) = \int_0^\infty [\beta_0]_7 a_7 e^{-\alpha_7 a_7} \rho_7(t,a_7)\, da_7$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + c_7 P_7 Q_7$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_7}{dt} = \omega_7 P_7 - \lambda_7 P_7 Q_7$$

$$\frac{dQ_7}{dt} = b_7 Q_7 - c_7 P_7 Q_7$$

Where $\omega_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$

## 33. THE EQUILIBRIUM $(\rho_0)_7(a_7), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_7(a_7), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_7 + (\rho_a)_7 = -\big(\mu_7(a_7) + (f_7)_7(P_7)\big)\rho,$$

$$\rho_7(t,0)$$

$$= \int_0^\infty \beta_7(a_7)\rho_7(t,a_7)da_7 \qquad (13)$$

The solution of (13) has the form of

$$\rho_7(t,a_7) = Ce_7^{r(t-a_7)} - \int_0^a (\tau)d\tau - \int_0^7 (f_7)_7\big(P_7(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_7(t) = C_7 e^{r_7 t - \int_0^t (f_7)_7(P_7(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_7^{-7}(r)$ if $k_7 > r$. This is equivalent to

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{(k_7)_7 a_7} da_7 < 1.$$

In fact, if we define

$$R_7(s_7) = \int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau)d\tau} e^{-s_7 a_7}\, da_7$$

$R(s)$ is a decreasing function. $R_7(r_7) = 1$ And $R_7(k_7)_7 < 1$ imply $r_7 < (k_7)_7$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_7(a_7), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{k_7 a_7} da_7$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_7 \equiv \int_0^\infty (\rho_0)_{77}(a_7)da_7$ *satisfies*

$$\int_0^{\infty} \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{-(f_7)_7(P_0)a_7} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_7(a_7), 0)$

Let $u_7(t, a_7) = \rho(t, a_7) - \rho_0(a_7)$ Then

$$N_7(t) \equiv \int_0^{\infty} u_7(t, a_7) da_7 = \int_0^{\infty} \rho_7(t, a_7) da_7 - \int_0^{\infty} (\rho_0)_7(a_7) da = P_7(t) - P_0)_7$$

And

$$(u_7)_7 + (u_a)_7 = (\rho_t)_7 + (\rho_a)_7 - (\rho_0^7)_7(a_7)$$

$$= -\mu_7(a_7) + (f_7)_7(P_7) + (f_7)_7(Q_7))(u_7 + (\rho_0)_7) + \mu_7(a_7)(\rho_0)_7 + (f_7)_7(P_0)_7(\rho_0)_7$$

$$= -\mu_7(a_7)u_7 - (f_7)_7(P_7) + (f_7)_7(Q_7))\, u_7 - (\rho_0)_7\, (f_7)_7(P_7) - (f_7)_7(P_0)_7$$

$$+ (f_7)_7(Q_7)), \qquad (15)$$

Where $(\rho_0^7)_7(a_7)$ denotes $d(\rho_0)_7(a_7)/da_7$

Expansion $(f_7)_7(P_7)$ and $(f_7)_7(Q_7)$ at $(P_0)_7$ and 0 yield respectively

$$(f_7)_7(P_7) = (f_7)_7((P_0)_7 + N_7)$$

$$= (f_7)_7(P_0)_7 + (f_7^7)_7(P_0)_7 N_7 + o_7(N_7), \qquad (16)$$

$$(f_7)_7(Q_7)$$

$$= (f_7^7)_7(0)Q + o_7(Q_7) \qquad (17)$$

Write $f_7(P_0) = f_{70}, f_7'(P_0) = f_{70}'$, and $f_7'(0) = f_{70}'$ and substitute (16) and (17) into (15) to obtain

$$(u_7)_7 + (u_a)_7 \approx -(\mu_7(a_7) + (f_{70})_7)u - (f_{70}^7)_7(\rho_0)_7(N_7) - (f_{70}^7)_7(\rho_0)_7(Q_7)$$

$$u_7(t, 0) = \rho_7(t, 0)(\rho_0)_7(0)$$

$$= h_7(Q_7)\int_0^{\infty} \beta_7(a_7)(u_7(t,a) + (\rho_0)_7(a_7)da_7 - (\rho_0)_7\,(0)$$

$$= h_7(Q_7)\int_0^{\infty} \beta_7(a_7)(u_7(t,a_7)da_7 + (h_7(Q_7) - 1)(\rho_0)_7\,(0) \qquad (18)$$

Expanding

$$h_7(Q_7) = 1 + \ (h^7)_7(0)Q_7 + 0(Q_7) \equiv 1 + h_0^{'}Q_7 + 0Q_7$$

And substituting it into (18), we arrive at

$$u_7(t,0) \approx \int_0^{\infty} \beta_7\,(a_7)u_7(t,a_7)da_7 + (h_{a_7}^{'})(\rho_0)_7Q_7$$





The linearization is the following :

$$[u_7] + [u_7] = -(\mu_7(a_7) + [t_{70}]_7u_7 - [\alpha_7]_7(a_7)N_7 - [\alpha_7]_7(a_7)Q_7$$

$$u_7(t,0) = \int_0^{\infty} \beta_7(a_7)u_7(t,a_7)da_7 + [\xi_7]Q_7 \qquad (19)$$

$$\frac{dQ_7}{dt} = -b_7Q_7$$

where

$$[\alpha_7]_7(a_7) = [t_7^7a]_7[p,o]_7(a_7)d2(a_7) = [f_{70}^7]_7(\rho\ o\ )_7(a_7)\xi = [\mathrm{h\ o}^7]_7[\rho\mathrm{o}]_7(\mathrm{o}) \text{ and } \mathrm{b}_7 = \mathrm{b}_7 - g_7(P_7)$$

If $g_7(P_0) > b$, clearly, $(\rho_0)_7(a_7), 0)$ is unstable and the predator population is persistent. Assume $g_7(P_0)_7 < b$.

Then, $Q_7(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_7 + u_a = -(\mu_7(a_7) + (f_{70})_7)u_7 - (\alpha_7)_7(a_7)N_7$$

$$= -(\mu_7(a_7) + f_{70})u_7 - \alpha_7(a_7)\int_0^\infty u_7(t, a_7)da_7$$

$$u_7(t, 0) = \int_0^\infty \beta_7\,(a_7)u_7(t, a_7)da_7$$

Using the method developed by Cushing [2, 3], we suppose

$$u_7(t, a_7) = \tilde{u}_7(a_7)^{e^{\gamma(t-a_7)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_7}\,da_7$. Then,

A solution of the first differential equation for $\tilde{u}(a_7)$ gives,

$$\tilde{u}_7(a_7) = e^{-\int_0^a \mu_7(\tau_7)d\tau - (f_{70})_7 a}\left(\tilde{u}_7(0) - w\int_0^a e^{\int_0^\tau \mu_7(s_7)ds + (f_{70})}\tau_7\,(\alpha_7)_7(\tau_7)e^{\gamma_7\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_7(a_7)e^{-\gamma_7 a_7 - (f_{70})_7 a_7 - \int_0^a \mu_7(\tau_7)d\tau}\right)\tilde{u}_7(0)$$

$$+w\int_0^\infty \beta_7(a_7)e^{-\gamma_7 a_7 - (f_{70})_7 a_7 - \int_0^a \mu_7(\tau_7)d\tau}\int_0^a (\alpha_7)_7(\tau_7)e^{-\gamma_7\tau_7}e^{(f_{70})_7\tau_7 + \int_0^\tau \mu_7(s_7)ds}d\tau_7\,da_7$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty} e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau} da\right)(\tilde{u}_7(0))_7$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau}\int_0^a e^{(f_{70})_7\tau_7+\int_0^\tau \mu_7(s_7)ds+y\tau}\alpha_7(\tau)d\tau\, da\right)w=0. \quad (23)$$

Define $Y(\gamma,a)=e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(t_7)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)da_7\right)u_7(0)$$

$$+w\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\, da_7=0, \quad (24)$$

$$\left(\int_0^{\infty}Y_7(\gamma_7,a_7)da_7\right)\tilde{u}_7(0)$$

$$+\left(-1-\int_0^{\infty}Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\, da_7\right)w=0.$$

Equation (20) has a solution of the form of $u_7(t,a_7)=u_7(a_7)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_7(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_7(\gamma_7,a_7)=\int_0^a(\rho_0)_7(\tau_7)Y^{-7}(\gamma_7,a_7)d\tau_7$. We obtain the characteristic equation

$$f_{70_7}^{l}\left(\int_0^{\infty} Y_7(\gamma_7,a_7)da_7\right)\left(\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$+\left(1-\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)da\right)\left(1+(f_{70}^{l})_7\int_0^{\infty}Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_7(a_7)=(\rho_0)_7(0)e^{-\int_0^a \mu_7(\tau_7)d\tau_7-(f_{70})_7a_7}$$

$$I_7(\gamma_7,a_7)=\int_0^a(\rho_0)_7(\tau_7)e^{\gamma_7\tau_7+(f_{70})_7\tau+\int_0^{\tau}\mu_7(s_7)ds_7}d\tau$$

$$=\frac{(\rho_0)_7(0)}{\gamma_7}(e^{\gamma_7a_7}-1)$$

Let

$$E_7(a_7)=e^{-(f_{70})_7a_7-\int_0^a\mu_7(\tau_7)d\tau_7}$$

Then $[Y_7]_7[\gamma_7]_7,a_7)=E_7(a_7)e^{-\gamma_7a_7}$. by a simple calculation,

$$\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da=\frac{[\rho_0]_7(0)}{\gamma}\left(1-\int_0^{\infty}\beta_7(a_7)E_7(a_7)e^{-\gamma_7a_7}\ da\right);$$

$$\int_0^{\infty}Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da=\frac{[\rho_0]_7(0)}{\gamma_7}\left(\int_0^{\infty}E_7(a_7)e^{-\gamma_7a_7}\ da_7\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_7(0)}{\gamma_7}[f_{70}^7]\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\,(\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7)$$

$$+\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right).$$

$$\left(1+\frac{(\rho_0)_7(0)}{\gamma_7}(f_{70}^{'})_7\left(\int_0^\infty E_7(a_7)da_7-\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right)(1+\frac{(f_{70}^{'})_7}{\gamma_7}(\rho_0)_7(0)\int_0^\infty E_7(a_7)da_7)=0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)da_7 = 1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7 \neq 0,$$

Since $\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)\int_0^a \alpha_7(\tau_7)Y^{-7}(\gamma_7,\tau_7)d\tau_7 da_7 > 0$. Therefore, the characteristic solution is

$$\gamma_7 = -(f_{70}^{'})_7(\rho_0)_7\int_0^\infty E_7(a_7)da_7.$$

Since $f_7'(P_0)_7 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_7), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE OF DEBIT AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_7), 0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_7 + (\mu_7(a_7) + [f_7]_7(P_7^*) + [f_7]_7(Q_7^*))P_7^* = 0, \quad (76)$$

$$p_7^*(0) = h_7(Q_7^*)\int_0^\infty \beta_7(a_7)p_7^*(a_J)p_7^*(a_7)da_7 \quad (77)$$

$$g_7(P_7^*) - b_7 = 0, \quad (78)$$

Where $P_7^* = \int_0^\infty p_7^*(a_7)\, da_7$

Equation (28) yields $P^* = g^{-7}(b)$,

From (26),

$$p_7^*(a) = p_7^*(0)e^{\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(Q_7^*)a_7}$$

Substituting into (27), we need to solve

$$\left(1 - h_7(Q_7^*)\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(Q_7^*)a_7}da_7\right)p_7^*(0) = 0$$

For $Q_7^*$

Fix $p_7^*(0)$ and $P_7^*$ Let

$$H_7(z_7) = 1 - h_7(z_7)\int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(z_7)a_7}da_{,7}$$

And assume that

$$[k_3]_7\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(g_7^{-7}(b_7))a_7 - [k_7]_7 a_7}\, da_7 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_7(z_7)] > 0.$$

Clearly,

$$\frac{dH_7}{dz_7} = h_7(z)(\,[f_7']_7(z_7)\,)\int_0^\infty a_7\beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - f_7(P_7^*)a_7 - f_7(z_7)a_7}da_{7}$$

$-h'(z_7)\int_0^\infty \beta_7(a_7)\, e^{\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(z_7)a_7}\quad da_7$

(See author's It is a debit-credit world)

By the assumptions on $\left[f_7^7\right]_7 and\ [h^7]^7 dH_7/dz > 0$. Hence, if

$$H_7(0) = 1 - \int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7} da_7 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_7(0) < 0$ follows if $g_7^-(b_7) < [f_7^{-7}]_7(r_7)$, where $r_7$ is the solution of (12).

Indeed, since

$$R_7(s_7) = \int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 e^{-y\frac{a}{7}}} da_7$$

Is a decreasing function of $s_7$ when $P_7^* = [g^{-7}]_7(b_7) < [f_7^{-7}]_7(r_7)$, that is $[f_7]_7(P_7^*) < r$, $R\left(f_7(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

**THEOREM 3**

Assume that the hypothesis (H3) holds and $R(f_7(g^{-7}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7\left(g_7^{-7}(b)\right)a_7} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_7^*(a_7), Q_7^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_7^*(a_7), Q_7^*)$ is more complicated. We again consider perturbations $m_7(t_7 a_7) = p_7(t_7, a_7) - p_7^*(a_7)$ and $W_7(t) = Q_7(t) - Q^*$.Then

$M_7(t) \equiv \int_0^\infty m_7(t,a_7)da_7 = P_7(t) - P_7^*,$

$$[m_7]_7 + [m_a]_7 = [p_7]_7 + [p_a]_7 - [p^*]_7(a_7) \approx -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7$$

$$-[f_{7*}^7]_7 p_7^*(a_7)M - [f_{7*}^7]_7 p_7^*(\mathrm{a})W_7,$$

$$m_7(t,0)p_7(t,0) - p_7^*(0) \approx h_* \int_0^\infty \beta_7(a_7)m_7(t,a_7)\,da + \frac{[h_*^7]_7}{[h_*]_7}\, p_7^*(0)W_7$$

$$\frac{dW_7}{dt} = \frac{dQ_7}{dt} \approx [g_*^7]_7 Q_7^* M,_7$$

Where $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}']_7$ denoting $[f'_7(P^*)]_7[f_{7*}']_7$ denoting

$[f_7]_7$ $(Q_7^*,$ $h_{7}^*$ denoting $h_*'$ denoting $h_7'$ $(Q_7^*)$, and $[g_*']_7$

Denoting $[g']_7(P_7^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$[m_7]_7 + [m_a]_7 = -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7 - [f'_{7*}]_7[p_7^*](a_7)M_7 - [f_{7*}']_7 p_7^*(a_7)W_7$

$m_7(t,0) = [h_*]_7 \int_0^\infty \beta_7(a_7)m_7(t,a_7)da_7 + \frac{[h_*^7]_7}{[h^*]_7}[p^*]_7(0)\,W,_7$ (30)

$$\frac{dW_7}{dt} = [g_*^7]_7[Q_7^*]M_7$$

Integrating $[m_7]_7(t,a_7)$ with respect to a from 0 to ∞, we have

$$\frac{dM_7(t)}{dt} = \int_0^\infty [mt]_7(t, a_7)da_7$$

$$= -\int_0^\infty [m_a]_7\,(t, a_7)da_7 - \int_0^\infty \mu(a_7)m(t, a_7)da_7$$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f'_{7*}]_7 \int_0^\infty p_7^*(a_7)da_7\right)M_7 - [f_{7*}^7]_7 \int_0^\infty p_7^*(a_7)da_7 W_7$$

$= m_7(t, 0) - \int_0^\infty \mu_7(a_7)\, m_7(t, a_7)da_7$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 - [f_{7*}^{'}]_7 P_7^* W_{.7}$$

From (30), we arrive at the equations

$$\frac{dW_7(t)}{dt} = \int_0^\infty ([h_*]_7\beta_7(a_7) - \mu_7(a_7))m_7(t, a_7)da_7$$

$$-\left([f_{7*}] + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 + \left(\frac{[h_*^{'}]_7}{[h_*]^7}\, P_7^*(0) - f_{7*}^{'} P_7^*\right)W_7,$$

(31)

$$\frac{dW_7(t)}{dt}[g'_*]_7\, Q_7^*\, M_{.7}$$

By the mean value theorem,

$$\int_0^\infty ([h_{"*}]_7\beta_7(a_7) - \mu_7(a_7))m_7(t, a_7)da_7 = ([h_*]_7\beta_7(\theta_7) - \mu_7(\theta_7))M,_7$$

Where $0< \theta_7 = \theta_7(t) < \infty$.

And if a $_7^7$ [a $_7^7$]2<4 [a $_7$]2<4 a $_7$e,

Re $\lambda_7[\lambda_7]_7 = \mathrm{Re}\lambda_7[\lambda_7]_7 = -\frac{7}{7}a_7[\mathrm{a}2\ ]_7 < 0$.

Let $\Delta$=max$_{7<u<\infty}|\mathrm{h}_*[\mathrm{h}_*]_7\beta_7(a_7) - \mu_7(a_7)|$. Then$|n_7(\mathrm{t})|$ 24∀t 0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium($\rho^*(\mathrm{a}), \boldsymbol{Q}^*$) is asymptotically stable if either a $_7^7$ <4 a $_7$ e and [a $_7$] $> 1\Delta$,or [a $_7^7$] 24[a $_7$]e and $(\mathrm{a}_7 - \sqrt{(\mathrm{a}_7^7)_7 - 4(\mathrm{a}_7)_7\mathrm{e}} > 1\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{\mathrm{a}}{\mathrm{a}+7}\right)$ and $\beta_7(a_7) = \mathrm{ae}^{-\mathrm{ma}}$ Then $\Delta$=[$\mu_0$] + 2 Simple calculation leads to

$$\int_0^\infty \beta_7(a_7)\mathrm{e}^{-\int_0^{\mathrm{a}}\mu_7(\tau_7)+d\tau_7\mathrm{da}} = \frac{1}{(\mathrm{m}+\mu_0+1)^7} + \frac{1}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_7$ + [$\mu_0$] 0.5. Assume[$f_7$]($x_7$)= $K_7x_7$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_7(x_7)$= [$0x_7$]/(2 + $x_7$); the functional responses [[$f_7$]($x$) = [0] $x_7/(1+x_7)$and $h_7(x_7)$=(1+$\omega_7x_7$)/(1 + $x_7$)(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_7 + [\mu_0] + [\mathrm{K}][P_0])}{[m_7]_7 + (\mu_0)_7 + [\mathrm{K}][P_0])_7^3\ [\mathrm{K}]_7[P_0]_7^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1+P_0) < b_7$ i e $[P_0] < b/(\delta_7 - b_7)$then the equilibrium$(P_0)$ $[P_0]_7(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_7$

$[f_{7*}] = K_7\ \mathrm{P}^*\ [f_{7*}] = \frac{\sigma \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*} \quad [h_*] = \frac{7+\omega_7 \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*}$

$[f_{7*}] = K_7\ [\mathrm{f}_{7*}^{\ 7}] = \frac{\sigma_7}{(7+\mathrm{Q}_7^*)^7} [\ \mathrm{g}_*^{\ 7}] = \frac{\sigma_7}{(7+\mathrm{P}_7^*)^7} \quad [\mathrm{h}_*^{\ 7}] = \frac{\omega_7 - 7}{(7+\mathrm{Q}_7^*)^7}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_7}{(m + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma)^7} + \frac{1\omega_7}{([m_7]_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma_7)^3} < 1$$

And

$$\frac{1}{(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^7} + \frac{1\omega_7}{(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $\mathrm{Q}^*$ which can be solved from the following equation:

$(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7} + (m_7 + [\mu_0]_7 + \frac{K_7 b_7}{\delta_7 - b_7} + \sigma_7)\ \mathrm{Q}_7^*)^3$

$= m_7 + [\mu_0]_7 + 3 + \frac{K_7 b_7}{\delta_7 - b_7}(1 + \mathrm{Q}^*)(1 + \omega_7 \mathrm{Q}_7^*$

Now

$$\left[\mathrm{a}\begin{matrix}1\\1\end{matrix}\right]-4\left[\mathrm{a}_{7}\right]\mathrm{e}=\left(\frac{1K_7 b_7}{\delta_7-b_7}\right)^{7}+\frac{4K_7 b_7\sigma_7\,\mathrm{Q}_{7}^{*}}{(\delta_7-b_7)(7+\mathrm{Q}_{7}^{*})}+\frac{\sigma_{7}^{7}\mathrm{Q}_{7}^{*7}}{\left(1+\mathrm{Q}_{7}^{*}\right)^{7}}$$

$$\frac{4\,b_7\sigma_7\,(\delta_{7}-\mathrm{b})\mathrm{Q}^{*}}{\delta_{7}\,(1+\mathrm{Q}_{7}^{*})^{7}}-\frac{4\,(\delta_{7}-b_7)^{7}\,(1-\omega_7)\mathrm{Q}_{7}^{*}}{\delta_{7}\left(1+\omega_7\mathrm{Q}_{7}^{*}\right)(1+\mathrm{Q}_{7}^{*})}\mathrm{P}*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_{7}\,ae_{7}^{mu}\;\;[\mu_0]_{7}-\frac{a_7}{a_7+7}$$

Hence if $[\mathrm{a}_{7}^{7}>4(a_7)_{7}\mathrm{e}$ and

In the following are given some more modular structures that are applied with the model and the concomitant results derived thereof:

Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments

(1) We have an intuitive understanding of how the world works. Our intuition works well in our everyday lives, but it doesn't quite match (e&eb) reality.

(2) One of the basic assumptions authors make about objects is that they are (=) objectively real.

(3) For example, when you leave your apartment in the morning, you imagine your house is right where you left it as you go about your day. Barring some calamity, you expect to find your apartment just as you left it when (e) you return in the evening.

(4) This idea is known as Macrorealism, in that a macroscopic object such as your house is still there even when (e) you're not looking at it

(5) In scientific terms we would say the state of an object exists independent of (e) an observer. While this seems obviously true, we can test the idea experimentally.

(6) To show this, let's consider an experiment. Suppose we have a box that contains a coin. At any time we can open the box and (eb) see whether it's heads or tails.

(7) The catch is (=) that the box has to follow a certain rule. Whatever you observe initially, (eb) you must observe the same thing ten seconds later.

(8) If you open the box and see heads, then (eb) ten seconds later when you open the box you must also see heads.

$$[a_7] = \frac{1K_7 b_7}{\delta_7 - b_7} + \frac{\sigma_7 \, \mathrm{Q}_7^*}{1 + \mathrm{Q}_7^*} > 1\vartriangle$$

This equation is important in its consequence in respect of following cases for which the model is applied:

**Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments**

(19) For the experiment we'll open the box, then open it five seconds later, then finally open it again at ten seconds. We know the first and last observations have to agree, but what about the observation in the middle? It depends upon what the box does with (e&eb) the coin.

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| (20) One possibility is that the box does nothing, so if you see tails the first time (eb), you will see tails both other times. |
| (21) Another possibility is that each time you open the box (eb) it flips the coin when you close it. |
| (22) So the first time you see heads, (eb) the second time tails, and (eb) the third time heads again. |
| (23) But they could also make a tricky box that flips (e&eb) the coin at seemingly random times, but does so in such a way that every ten seconds it repeats (eb) the pattern. In that case, authors would always see the same thing the first and last time, but the middle observation would seem (eb) random. |
| (24) If Macrorealism is true, then (eb) middle observations being always the same as (=) the first and last observation, always opposite or always random are (=) the only possibilities for the experiment |

**Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry Download as PDF (Size: 62KB) HTML PP. 50-53 DOI: 10.4236/jqis.2011.12007 5,508 Downloads 14,712 Views Citations Author(s) Leave a comment Mohamed S. El Naschie**

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| (19) Building upon the pioneering work of J. Bell [1] and an incredible result due to L. Hardy [2] it was shown (eb) that the probability of quantum entanglement of two particles is (=) a maximum of 9.0169945 percent [2]. |
| (20) This happens to be exactly the golden mean to (e) the power of five (?5) [3-7]. |
| (21) Although it has gone largely unnoticed for a long time, this result was essentially established independently in a much wider context by (e) the present author almost two decades ago [3-6]. |
| (22) The present work gives two fundamentally different derivations of (e) Hardy's beautiful result leading to (eb) precisely the same general conclusion, namely that by virtue of (e) the zero measure of the underlying |
| (23) Cantorian-fractal spacetime geometry the notion of (e) spatial separability in quantum physics is (=) devoid of any meaning [7]. The first derivation is purely logical and uses (e) a probability theory which combines the discrete with (e&eb) the continuum. |
| (24) The second derivation is purely geometrical and topological using (e) the fundamental equations of a theory developed by the author and his collaborators frequently referred to as E-infinity or Cantorian spacetime theory [3-7]. KEYWORDS Hardy's Quantum Entanglement, Golden mean, Cantor sets, Fractal Spacetime, E-Infinity Theory, Quantum Mechanics, J. S. Bell Cite this paper M. El Naschie, "Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry," Journal of Quantum Information Science, Vol. 1 No. 2, 2011, pp. 50-53. Doi: 10.4236/jqis.2011.12007. |

**Chaos, Solitons & Fractals Volume 28, Issue 4, May 2006, Pages 857–878 Sir Hermann Bondi 1919–2005 Cantorian spacetime and Hilbert space: Part I—Foundations G. Iovane http://dx.doi.org/10.1016/j.chaos.2005.08.074**

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| (31) Authors are going to show the link between the ε(∞) Cantorian space and( e&eb) the Hilbert spaces H(∞). |
| (32) In particular, El Naschie's ε (∞) is (=) a physical spacetime, i.e. an infinite dimensional fractal space, where (e) time is spacialized and the transfinite nature manifests (eb) itself. |
| (33) El Naschie's Cantorian spacetime is an arena where (e) the physics laws appear at each scale in a self-similar way linked to (e&eb) the resolution of the act of observation. |
| (34) By contrast the Hilbert space H (∞) is (=) a mathematical support, which describes (eb) the interaction between the observer and (e&eb) the dynamical system under (e&eb) measurement. |

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| (35) The present formulation, which is based on (e) the non-classical Cantorian geometry and topology of spacetime, automatically solves (e) the paradoxical outcome of (e) the two-slit experiment and the so-called particle–wave duality. |
| (36) In particular, measurement (i.e. the observation) is equivalent to (=) a projection of ε(∞) in the Hilbert space built on 3 + 1 Euclidean spacetime. |
| (37) Consequently, the wave–particle duality becomes (=) a mere natural consequence of (e) conducting an experiment in (eb) a spacetime with (e&eb) non-classical topological and (e&eb) geometrical structures, while (e) observing and taking measurements in (eb) a classical smooth 3 + 1 Euclidean spacetime. |
| (38) In other words, the experimental fact that a wave–particle duality exists is (=) an indirect confirmation of the existence of ε(∞) and a property of (e) the quantum–classical interface. |
| (39) Another direct consequence of the fact that real spacetime is) =) infinite dimensional hierarchical ε(∞) is (=) the existence of scaling law R(N), introduced by the author, which generalizes the Compton wavelength. |
| (40) It gives an answer to the problem of (e) segregation of matter at different scales, and shows (eb) the role of fundamental constants such as (=) the speed of light and Plank's constant h in the fundamental lengths scale without (e) invoking the principles of quantum mechanics. Copyright © 2005 Elsevier Ltd. All rights reserved |

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**Chaos, Solitons & Fractals Volume 14, Issue 9, December 2002, Pages 1341–1351 Fractal strings as an alternative justification for El Naschie's Cantorian spacetime and the fine structure constant Carlos Castro, http://dx.doi.org/10.1016/S0960-0779(02)00110-8**

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| (25) Beginning with the most general fractal strings/sprays construction recently expounded in the book by Lapidus and Frankenhuysen, it is shown how the complexified extension of El Naschie's cantorian-fractal spacetime model belongs to (e&eb) a very special class of families of fractal strings/sprays whose scaling ratios are given by (e) suitable pinary (pinary, p prime) powers of the Golden Mean. |
| (26) They then proceed to show why the logarithmic periodicity laws in (eb) Nature are direct physical consequences of (e) the complex dimensions associated with (e&eb) these fractal strings/sprays. |
| (27) Authors proceed with a discussion on quasi-crystals with (e&eb) p-adic internal symmetries, von Neumann's Continuous Geometry, the role of wild topology in (e&eb) fractal strings/sprays, (e&eb) the Banach-Tarski paradox, and tesselations of (e) the hyperbolic plane, quark confinement and (e&eb) the Mersenne-prime hierarchy of (e) bit-string physics in determining the fundamental physical constants in Nature. Fax: +1-914-965-3397 Copyright © 2002 Published by Elsevier Ltd. About ScienceDirectRemote access Shopping cart Contact and support Terms and conditions Privacy policy Cookies are used by this site. For more information, visit the cookies page. Copyright © 2016 Elsevier B.V. or its licensors or contributors. Science Direct ® is a registered trademark of Elsevier B.V. |

**Chaos, Solitons & Fractals Volume 21, Issue 2, July 2004, Pages 261–271 On Cantorian spacetime over number systems with division by zero Jakub Czajko, http://dx.doi.org/10.1016/j.chaos.2003.12.046**

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| (28) Division by zero is proposed in order to open (eb) new horizons for (e) abstract mathematics and physics. |
| (29) Although the formerly forbidden operation may be of (e) little practical value, the feasibility of unrestricted operations is needed for (e) meaningful deployment of linear vector spaces, whose duals are (=) inverse transposes of their primary vector spaces. |
| (30) Both quantum mechanics and (e&eb) pan-geometry of multispatial hyperspace operate on (e&eb) vectors that belong to (e&eb) mutually dual linear vector spaces. |
| (31) The division by zero also may remove (e) many years old paradox related to (e&eb) multiplication of zeros and makes it possible to introduce (eb) the concept of bigroup as (=) an additive and multiplicative group at the same time. |
| (32) This operational extension to (e&eb) algebraic groups makes thus (eb) an infinite-dimensional, fractal Cantorian spacetime (=) the preferred, invertible, operationally unrestricted, abstract mathematical infrastructure for physics. Tel.: +1-925-681-2058 Copyright © 2003 Elsevier Science Ltd. All rights reserved. |

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The nontrivial equilibrium is stable.

Some of the examples in which stability analysis is of seminal importance are: These are the system to which the

model becomes applicable

**Observation of a Dissipation-Induced Classical to Quantum Transition J. Raftery, D. Sadri, S. Schmidt, H. E. Türeci, and A. A. Houck Phys. Rev. X 4, 031043 – Published 8 September 2014**

(363)Here, authors report the experimental observation of (e) a dynamical quantum phase transition in (eb) a strongly interacting open photonic system.

(364) The system studied, comprising a Jaynes-Cummings dimer realized on (eb) a superconducting circuit platform, exhibits (eb) a dissipation-driven localization transition.

(365)Signatures of the transition in the homodyne signal and photon number reveal (eb) this transition to be from (e) a regime of classical oscillations into (e&eb) a macroscopically self-trapped state manifesting revivals, (=) a fundamentally quantum phenomenon.

(366) This experiment also demonstrates (eb) a small-scale realization of a new class of quantum simulator, whose well-controlled (e&eb) coherent and dissipative dynamics is suited to (e) the study of quantum many-body phenomena out of (e) equilibrium. Received 10 December 2013 DOI:http://dx.doi.org/10.1103/PhysRevX.4.031043 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**Viewpoint: Fermionic Vortices Find their Dual John McGreevy, Department of Physics, University of California at San Diego, La Jolla, CA 92093-0319, USA June 27, 2016• Physics 9, 72Theoretical work reveals a surprising relationship between the physics of fermionic vortices and quantum electrodynamics. Figure caption expand figure APS/Alan Stonebraker Figure 1: The action of discrete symmetries on the fermionic particle (top) and its dual fermionic vortex (bottom).**

(367)Dirac fermions in 2+1 dimensions (two spatial dimensions plus time) have been found in (eb) graphene and on (e&eb) the surfaces of topological insulators (TIs).

(368) A trio of papers now shows that these particles can also be found in (eb) a dual description of quantum electrodynamics (QED) [1–3].

(369)This deceptively simple statement joins (e&eb) together several disparate lines of progress in theoretical physics: interacting topological insulators, (e&eb) anomalies in quantum field theory, (e&eb) spin liquids, and (e&eb) compressible Quantum Hall States.

(370)Humans are not very good at understanding interacting many-body systems.

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| (371)We understand the behavior of free particles and can sometimes use (e) perturbation theory to (e) incorporate weak interactions between (e&eb) them. |
| (372)But what happens when we couple (e&eb) the particles strongly? |
| (373)To reveal the interesting collective phenomena that can result, we've had to rely largely on (e) experiments. |
| (374)Examples include (e) superconductivity (e&eb) superfluidity (e&eb) the integer and (e&eb) fractional quantum Hall effects (e&eb) the fractionalization of charge and spin, and (e&eb) the wetness of water. |
| (375)The fractional quantum Hall effect (FQHE) is a physical phenomenon in which the Hall conductance of 2D electrons shows (eb) precisely quantised plateaus at (eb) fractional values of (e) {\displaystyle e^{2}/h} e^{2}/h. |
| (376) It is a property of (e) a collective state in which electrons bind (e&eb) magnetic flux lines to make (eb) new quasiparticles, and excitations have (e) a fractional elementary charge and possibly also fractional statistics. The 1998 Nobel Prize in Physics was awarded to Robert Laughlin, Horst Störmer, and Daniel Tsui for the discovery and explanation of the fractional Hall effect.[1] |
| (377) However, Laughlin's explanation was (=) a phenomenological guess and only applies to (e&eb) fillings {\displaystyle \nu =1/m} \nu =1/m where (e) {\displaystyle m} m is an odd integer. The microscopic origin of the FQHE remains unknown and is a major research topic in condensed matter physics. (Wikipedia) |
| (378)The quantum Hall effect (or integer quantum Hall effect) is a quantum-mechanical version of (e) the Hall effect, observed in (eb) two-dimensional electron systems subjected to (e&eb) low temperatures and strong magnetic fields, in which the Hall conductance σ undergoes (e&eb) certain quantum Hall transitions to take on (e&eb) the quantized values {\displaystyle \sigma ={\frac {I_{\text{channel}}}{V_{\text{Hall}}}}=\nu \;{\frac {e^{2}}{h}},} \sigma = \frac{I_\text{channel}}{V_\text{Hall}} = \nu \; \frac{e^2}{h}, where {\displaystyle I_{\text{channel}}} I_\text{channel} is the channel current, {\displaystyle V_{\text{Hall}}} V_\text{Hall} is the Hall voltage, e is the elementary charge and h is Planck's constant. The prefactor ν is known as the "filling factor", and can take on either integer (ν = 1, 2, 3, ...) or fractional (ν = 1/3, 2/5, 3/7, 2/3, 3/5, 1/5, 2/9, 3/13, 5/2, 12/5, ...) values. |
| (379) The quantum Hall Effect is referred to as (=) the integer or fractional quantum Hall effect depending on whether ν is an integer or fraction, respectively. |
| (380)The striking feature of the integer quantum Hall Effect is (=) the persistence of the quantization (i.e. the Hall plateau) as (e) the electron density is varied. |
| (381) Since the latter remains constant when (e) the Fermi level is in a clean spectral gap, it must be that the Fermi level is immersed in (eb) essential spectrum (i.e. of finite density of states), though Anderson localized. |
| (382) The quantization and stability of the Hall conductance in such conditions was achieved in 1994 by (e) Bellissard, van Elst and Schulz-Baldes using the methods of (e) Non-Commutative Geometry.[1] |
| (383)The fractional quantum Hall effect is more complicated, as its existence relies fundamentally on (e) electron–electron interactions. |
| (384)Although the microscopic origins of the fractional quantum Hall effect are unknown, there are several phenomenological approaches that provide (eb) accurate approximations. |
| (385) For example the effect can be thought of as (=) an integer quantum Hall effect, not of electrons but of (e) charge-flux composites known as (=) composite fermions. |

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| (386) In 1988, it was proposed that there was quantum Hall effect without Landau levels. This quantum Hall Effect is referred to as (=) the quantum anomalous Hall (QAH) effect. |
| (387)There is also a new concept of the quantum spin Hall effect which is an analogue of the quantum Hall effect, where (e) spin currents flow instead of (e&eb) charge currents.[2] |
| (388)In two dimensions, when classical electrons are subjected to (e&eb) a magnetic field they follow (eb) circular cyclotron orbits. |
| (389)When the system is treated quantum mechanically (eb) these orbits are quantized. |
| (390)The energy levels of these quantized orbitals take on (e) discrete values |
| (391):{\displaystyle \textstyle E_{n}=\hbar \omega _{c}(n+1/2),} \textstyle E_{n}=\hbar \omega _{c}(n+1/2), where ωc = eB/m is the cyclotron frequency |
| (392) These orbitals are known as Landau levels, and at weak magnetic fields, their existence gives rise to (eb) many interesting "quantum oscillations" such as the Shubnikov–de Haas oscillations and the de Haas–van Alphen effect (which is often used to map the Fermi surface of metals). |
| (393) For strong magnetic fields, each Landau level is (=) highly degenerate (i.e. there are many single particle states which have the same energy En). |
| (394)Specifically, for a sample of area A, in magnetic field B, the degeneracy of each Landau level is (=) {\displaystyle \textstyle N=g_{s} BA/\phi _ {0},} \textstyle N=g_{s} BA/\phi _ {0}, where gs represents (eb) a factor of 2 for spin degeneracy, and ϕ0 = 2·10−15 Wb is the magnetic flux quantum. |
| (395) For sufficiently strong B-fields, each Landau level may have (e) so many states that all of the free electrons in the system sit in only a few Landau levels; it is in this regime where (e) one observes the quantum Hall Effect. |
| (396)The integers that appear in the Hall Effect are (=) examples of **topological quantum numbers**. |
| (397) They are known in mathematics as the first **Chern numbers** and are closely related to (e&eb) **Berry's phase.** |
| (398)A striking model of much interest in this context is the Azbel-Harper-Hofstadter model whose quantum phase diagram is (=) the Hofstadter butterfly shown in the figure. (deleted due to spatial constraints) |
| (399)The vertical axis is (=) the strength of the magnetic field and the horizontal axis is the chemical potential, which fixes the electron density. The colors represent the integer Hall conductance. Warm colors represent positive integers and cold colors negative integers. The phase diagram is fractal and has structure on all scales. In the figure there is an obvious self-similarity. |
| (400)Concerning physical mechanisms, impurities and/or particular states (e.g., edge currents) are important for (e) both the 'integer' and 'fractional' effects. |
| (401) In addition, Coulomb interaction is also essential in (eb) the fractional quantum Hall effect. |
| (402)The observed strong similarity between integer and (=) fractional quantum Hall effects is explained by (e) the tendency of electrons to form bound states with (e&eb) an even number of magnetic flux quanta, called composite fermions. |
| (403)Duality is the circumstance in which cranking up interactions results in (eb) another system we understand, usually one that can be described in terms of (e&eb) another set of weakly coupled particles. |
| (404)When such a dual representation is found, (eb) it gives a solution to the strong coupling problem. |
| (405) A further reason that theoretical physicists get excited about dualities is that they undermine (e) the very notion of an “elementary particle.” |
| (406)In terms of the original variables, the dual, weakly coupled particles are (=) generally solitons—large, heavy, classical objects made up of (e) many quanta. |

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| (407)The remarkable thing is that solitons can become (=) "elementary," in the sense that the model that describes the system can be reformulated with the solitons as (=) the basic constituents. |
| (408)An example of a duality, which we will use below, arises in (eb) Maxwell's theory of electromagnetism, where (e) the strength of the interactions is determined by (e) the electric charge e. |
| (409)The duality exchanges electric and (e&eb) magnetic point charges, and the interaction strength in the dual description is 1⁄e, which is small when (e) e is large. |
| (410)This idea has been generalized to (e) many other field theories, thanks, in particular, to tools of (e) supersymmetry [4]. |
| (411)Another classic example of duality involves (e&eb) vortices. |
| (412) A vortex is a localized object in the plane around which a phase field winds (e&eb). (A phase field takes values on the circle.) |
| (413)Dualities relating particles and vortices have (e) a rich history in particle physics and condensed-matter physics. |
| (414) The simplest version relates the particles in the normal state of a superfluid to (e&eb) the vortices in a dual superconductor [5]. |
| (415)This charge-vortex duality has been (=) a workhorse in the study of two-dimensional condensed-matter systems, such as (=) fractional Quantum Hall States, superconducting thin films, and beyond-Landau critical phenomena [6]. |
| (416)Vortices can be (=) fermions (as they frequently are, for instance, in type-II superconductors). |
| (417) A dual description of a system with (e&eb) fermionic vortices thus requires (e) a new ingredient, since (e) it must include (e) such a fermionic particle. |
| (418)The new work [1–3] builds on (e&eb) earlier theoretical work [7] to make (eb) an appealing proposal for (e) the required duality: the fermionic vortex of QED in (eb) 2+1 dimensions is (=) a free Dirac fermion. The argument for this proposal builds on our understanding of TIs. |
| (419)TIs are distinguished only by (e) their edge states. |
| (420) The most famous example is a (3+1)-dimensional time-reversal invariant TI, which can host boundary states whose band structure is defined by (e) a single Dirac cone. |
| (421)A single time-reversal-symmetric Dirac cone is not possible without (e) the bulk TI, so its appearance in experiments provides (eb) a signature of the bulk phase. |
| (422) But the Dirac cone is not (e) the only possible signature. |
| (423) Deformations of the boundary conditions can dramatically change (e&eb) the spectrum of the edge states, yet preserve some essential "TI-ness." This essential property is (=) called an anomaly in high-energy theory. |
| (424)Max Metlitski of the University of California, Santa Barbara, and Ashvin Vishwanath of the University of California, Berkeley [1], and independently, Chong Wang and Todadri Senthil at the Massachusetts Institute of Technology, Cambridge [2], use this ambiguity in the edge theory and the electric-magnetic duality of an auxiliary theory of bulk electromagnetism coupled to (e&eb) the TI to argue for the duality. |
| (425)When I said that (3+1)-dimensional QED had a duality that interchanged the electric and magnetic charges, one could have (e) complained that the electron is (=) a fermion, while (e) a magnetic monopole may not be. |
| (426)However, the authors observe that in a TI, a charge-two monopole is (=) also a fermion. |

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| (427) By analyzing the spectrum of electric and magnetic charges, they conclude that the duality exchanges the electron and (e&eb) the charge-two monopole and, moreover, exchanges a TI with (e&eb) time-reversal symmetry with (e&eb) a "chiral TI," which instead has (e) particle-hole symmetry |
| (428). But this is just what is required of (e) a symmetry that exchanges a fermionic particle with (e&eb) a fermionic vortex: time-reversal symmetry maps the electron to itself but maps the vortex to (e&eb) another vortex whose phase winds (e&eb) the other way (see Fig. 1). |
| (429)The exchange of particles and holes, in turn, preserves (eb) the vortex. |
| (430)The bulk duality implies (eb) a duality between edge theories. |
| (431)The freedom to deform the edge of the TI implies (eb) a weak form of equivalence: although the edge theories are (=) topologically the same (in terms of symmetry assignments and (e&eb) anomalies), they need not have (e) identical low-energy physics. |
| (432)One side of the duality is (=) a free, massless Dirac particle, but the long-distance physics of (2+1)-dimensional QED is (=) a long-standing problem, which is important in the study of spin liquids. |
| (433) Many gauge theories in two dimensions develop (eb) a gap in their energy spectrum because of (e) the proliferation of monopoles. |
| (434) In the realization proposed by Metlitski and Vishwanath [1] and by Wang and Senthil [2], this cannot happen because (e) the monopoles carry electric charge, and indeed the stronger notion of duality as (=) an equivalence of low-energy physics may hold. |
| (435)Evidence for the stronger relation comes from (e) David Mross and colleagues at the California Institute of Technology, Pasadena [3], who have derived a lattice version of the duality using (e) a clever but explicit change of variables. |
| (436)Their construction is particularly interesting because (e) it uses (e) a lattice of one-dimensional models which themselves can only occur at the edge of a TI. |
| (437)Authors explicitly derive the duality between a free electronic Dirac cone and (e&eb) quantum electrodynamics in (eb) (2+1) dimensions (QED3) with (e&eb) N=1 fermion flavors. |
| (438)The duality proceeds via (e&eb) an exact, nonlocal mapping from electrons to (e&eb) dual fermions with (e&eb) long-range interactions encoded by (e) an emergent gauge field. |
| (439)This mapping allows (eb) us to construct (eb) parent Hamiltonians for exotic topological-insulator surface phases, derive (eb) the particle-hole-symmetric field theory of (e) a half-filled Landau level, and nontrivially constrain (e) QED3 scaling dimensions. |
| (440)They similarly establish duality between bosonic topological insulator surfaces and (e&eb) N=2 QED3. |

**Black hole thermodynamics with generalized uncertainty principle Li Xiang, X. Q. Wen (Submitted on 6 Jan 2009 (v1), last revised 18 Mar 2009 (this version, v2))**

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| (441) In the standard viewpoint, the temperature of a stationary black hole is proportional to (=) its surface gravity, $T_H=\hbar\kappa/2\pi$. |
| (442)This is a semiclassical result and the quantum gravity effects are not (e) taken into consideration. |
| (443)This Letter explores a unified expression for (e) the black hole temperature in the sense of (e) a generalized uncertainty principle (GUP). |
| (444) Discussion involves a heuristic analysis of (e) a particle which is absorbed by (e) the black hole. |
| (445)Besides a class of static and spherically symmetric black holes, an axially symmetric Kerr-Newman black hole is considered. Different from the existing literature, authors suggest (eb) that the black hole's irreducible mass represent (eb) the characteristic size in (eb) the absorption process. |

(446)The information capacity of a remnant is also discussed by (e) Bousso's D-bound in de Sitter spacetime. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP 0910:046,2009 DOI: 10.1088/1126-6708/2009/10/046 Cite as: arXiv:0901.0603 [gr-qc] (or arXiv:0901.0603v2 [gr-qc] for this version)

IF we define

$$[a_7]_7 = [f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}]\,P_7^*$$

$$[a_7]_7 = \left[f_{7*}^{7}\right] P_7^* - \frac{h_*^7}{h_*} P^*(0)$$

$$e = [g_*^7]_7\, Q_7^*$$

$$\eta_7(t) = \left[h_*\right]_7 \beta_7 \left(\theta_7(t)\right) - \mu_7(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\,[a_7]M_7 - [a_7]_7 W_7 - [n(t)]M_7$$

$$w_7(t) = eM_7$$

Write

$$A = \begin{pmatrix} -[a_7]_7 - [a_7]_7 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)= \begin{pmatrix} [\eta(t)]_7 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_7 \\ W_7 \end{pmatrix}$$

The solution of (32)is

$$X_7(t) = e^{A,t}[X]_7(0) + \int_0 e^{At}$$

## PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION

## ATTRITION MODELS

## MODULE EIGHT

Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal.

### 34. MODEL FORMULATION

**Model is applied to systems given in the following:**

**Researchers blur the line between classical and quantum physics by connecting chaos and entanglement July 12, 2016 by Sonia Fernandez**

(2078) Experimental link between quantum entanglement (left) and (e&eb) classical chaos (right) found using (e) a small quantum computer. Credit: University of California - Santa Barbara Using a small

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| quantum system consisting of three superconducting qubits, researchers at UC Santa Barbara and Google have uncovered a link between aspects of classical and (e&eb) quantum physics thought to be unrelated: classical chaos and (e&eb) quantum entanglement. |
| (2079) Their findings suggest (eb) that it would be possible to use (e) controllable quantum systems to investigate certain fundamental aspects of nature. |
| (2080) "It's kind of surprising because (e) chaos is this totally classical concept—there's (=) no idea of chaos in a quantum system," Charles Neill, a researcher in the UCSB Department of Physics and lead author of a paper that appears in Nature Physics. |
| (2081) "Similarly, there's no concept of entanglement within (eb) classical systems. And yet it turns out that chaos and (e&eb) entanglement are really very strongly and clearly related." |
| (2082) Initiated in the 15th century, classical physics generally examines and describes (eb) systems larger than atoms and molecules. |
| (2083) It consists of hundreds of years' worth of study including Newton's laws of motion, electrodynamics, relativity, thermodynamics as well as chaos theory—the field that studies the behavior of (e) highly sensitive and unpredictable systems. |
| (2084) One classic example of chaos theory is the weather, in which a relatively small change in one part of (e) the system is enough to foil predictions—and vacation plans—anywhere on the globe. |
| (2085) At smaller size and length scales in nature, however, such as those involving (e&eb) atoms and photons and their behaviors, classical physics falls short. |
| (2086) In the early 20th century quantum physics emerged, with its seemingly counterintuitive and sometimes controversial science, including (e) the notions of superposition (the theory that a particle can be located in (eb) several places at once) and entanglement (particles that are deeply linked behave as (=) such despite physical distance from one another). |
| (2087) And so began the continuing search for connections between the two fields. All systems are fundamentally quantum systems, according Neill, but the means of describing in a quantum sense the chaotic behavior of, (e) say, air molecules in an evacuated room, remains limited. |
| (2088) Imagine taking a balloon full of air molecules, somehow tagging them so you could see them and then releasing them (eb) into a room with no air molecules, noted co-author and UCSB/Google researcher Pedram Roushan. |
| (2089) One possible outcome is that the air molecules remain clumped together (e&eb) in a little cloud following the same trajectory around the room |
| (2090) And yet, he continued, as we can probably intuit, the molecules will more likely take off in (eb) a variety of velocities and directions, bouncing off walls and interacting with (e&eb) each other, resting after the room is (=) sufficiently saturated with them. |
| (2091) "The underlying physics is (=) chaos, essentially," he said. |
| (2092) The molecules coming to rest—at least on (eb) the macroscopic level—is the result of (e) thermalization, or of reaching equilibrium after they have achieved (eb) uniform saturation within the system. |
| (2093) But in the infinitesimal world of quantum physics, there is (=) still little to describe that behavior. |
| (2094) The mathematics of quantum mechanics, Roushan said, do not allow for (e) the chaos described by Newtonian laws of motion. |
| (2095) To investigate, the researchers devised an experiment using (e) three quantum bits, the basic computational units of (e) the quantum computer. |

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| (2096) Unlike classical computer bits, which utilize (e) a binary system of two possible states (e.g., zero/one), a qubit can also use (e) a superposition of both states (zero and one) as a single state. |
| (2097) Additionally, multiple qubits can entangle, or link so closely that their measurements will (=) automatically correlate. |
| (2098) By manipulating these qubits with electronic pulses, Neill caused them to interact, rotate and evolve (e&eb, e&eb, eb) in the quantum analog of (e) a highly sensitive classical system. |
| (2099) The result is a map of entanglement entropy of (e) a qubit that, over time, comes to strongly resemble that of (e) classical dynamics—the regions of entanglement in (eb) the quantum map resemble the regions of (e) chaos on the classical map. |
| (2100) The islands of low entanglement in (eb) the quantum map are located in the places of (e) low chaos on the classical map. |
| (2101) "There's a very clear connection between entanglement and (e&eb) chaos in these two pictures," said Neill. |
| (2102) "And, it turns out that thermalization is (=) the thing that connects chaos and (e&eb) entanglement. |
| (2103) It turns out that they are actually the driving forces behind (eb) thermalization. |
| (2104) "What we realize is that in almost any quantum system, including on quantum computers, if you just let it evolve and you start to study what happens as a function of time, (eb) it's going to thermalize," added Neill, referring to the quantum-level equilibration. "And this really ties together the intuition between classical thermalization and (e&eb) chaos and how it occurs (eb) in quantum systems that entangle." |
| (2105) The study's findings have fundamental implications for (e) quantum computing. |
| (2106) At the level of three qubits, the computation is relatively simple, said Roushan, but as researchers push to build increasingly sophisticated and powerful quantum computers that incorporate (e) more qubits to study highly complex problems that are beyond the ability of classical computing—such as those in the realms of machine learning, artificial intelligence, fluid dynamics or chemistry—a quantum processor optimized for such calculations will be a very powerful tool. |
| (2107) "It means we can study things that are completely impossible to study right now, once we get to bigger systems," said Neill.More information: Nature Physics, DOI: 10.1038/nphys3830 , http://www.nature.com/nphys/journal/vaop/ncurrent/full/nphys3830.html Journal reference: Nature Physics search and more info website Provided by: University of California - Santa Barbara |

**Title: coherent control of macroscopic quantum states in a single-Cooper-pair box Authors: Nakamura, Y.; Pashkin, Yu. A.; Tsai, J. S. Publication: Nature, Volume 398, Issue 6730, pp. 786-788 (1999). (Nature Homepage) Publication Date: 04/1999 Origin: NATURE DOI: 10.1038/19718 Bibliographic Code: 1999Natur.398...786N**

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| (2108) A nanometre-scale superconducting electrode connected to (e&eb) a reservoir via (e&eb) a Josephson junction constitutes (e) an artificial two-level electronic system: a single-Cooper-pair box. |
| (2109) The two levels consist of (e) charge states (differing by 2e, where (e) e is the electronic charge) that are coupled by (e) tunnelling of Cooper pairs through (e&eb) the junction. |
| (2110) Although the two-level system is macroscopic, containing (e) a large number of electrons, the two charge states can be (=) coherently superposed. |

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| (2111) | The Cooper-pair box has therefore been suggested as (=) a candidate for a quantum bit or `qubit'- the basic component of a quantum computer. |
| (2112) | Here authors report the observation of quantum oscillations in (eb) a single-Cooper-pair box. |
| (2113) | By applying a short voltage pulse via (e&eb) a gate electrode, we can control (e&eb) the coherent quantum state evolution: the pulse modifies (e&eb) the energies of (e) the two charge states non-adiabatically, bringing them into (e&eb) resonance. |
| (2114) | The resulting state-a superposition of (e) the two charge states-is detected by (e) a tunnelling current through (e&eb) a probe junction. |
| (2115) | Results demonstrate electrical coherent control of (e) a qubit in a solid-state electronic device. |

**Superconducting quantum computing From Wikipedia, the free encyclopedia**

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| (2116) | Superconducting quantum computing is a promising implementation of (e) quantum information technology that involves (e&eb) nanofabricated superconducting electrodes coupled through (e&eb) Josephson junctions. |
| (2117) | As in a superconducting electrode, the phase and (e&eb) the charge are conjugate variables. |
| (2118) | There exist three families of superconducting qubits, depending on (=) whether the charge, (e&eb) the phase, or neither of (e) the two is good quantum number. |
| (2119) | These are respectively termed charge qubits (e&eb) flux qubits, and (e&eb) hybrid qubits. |

**Theory**

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| (2120) | Unlike many other physical implementations of a qubit which involve (e&eb) exclusively two-level systems (such as nuclear spin and (e&eb) photon polarization), the integrated quantum circuit involved in (e&eb) a superconducting qubit is (=) a multi-level system, of which only the first two levels are used as (=) the computational basis. |
| (2121) | A basic requirement for such an implementation is that the energy levels are (=) not uniformly spaced, so that (e) photons of a particular frequency which cause (eb) transition between the 0 and 1 levels do not(e) cause transitions from the first level to (e&eb) the higher levels as well. |
| (2122) | For electronic signals to be carried from one part of the circuit to (e&eb) another without energy loss and (e&eb) hence decoherence, the system also needs to be (=) non-dissipative, i.e., the metallic parts involved should have (e) zero resistance. |
| (2123) | These circuits currently need to be (=) operated at very low temperatures so that, 1) superconductivity is (=) realized, and 2) thermal fluctuations do not (e) cause transitions between (e&eb) energy levels. |

**Quantum LC Circuit**

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| (2124) | The simplest non-dissipative quantum circuit consists simply of (e) an inductor L {\displaystyle L} L and (e&eb) capacitor C {\displaystyle C} C, with (e&eb) the metallic wires connecting them being superconducting. |
| (2125) | Here the flux Φ {\displaystyle \Phi } \Phi through (e&eb) the inductor and (e&eb) the charge Q {\displaystyle Q} Q in the capacitor are (=) canonically conjugate variables which obey (eb) the commutation relation [ Φ , Q ] = i ℏ {\displaystyle [\Phi ,Q]=i\hbar } [\Phi ,Q]=i\hbar |

**The Hamiltonian associated with this system is**

(2126) H = Φ 2 2 L + Q 2 2 C {\displaystyle H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}} H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}

(2127) This gives rise to the energy levels E = ℏ ω 0 ( n + 1 2 ) {\displaystyle E=\hbar \omega _{0}(n+{\frac {1}{2}})} E=\hbar \omega _{0}(n+{\frac {1}{2}}) Here ω 0 = 1 L C {\displaystyle \omega _{0}={\frac {1}{\sqrt {LC}}}} \omega _{0}={\frac {1}{{\sqrt {LC}}}}

**Quantum gravity Carlo Rovelli (2008), Scholarpedia, 3(5):7117. doi:10.4249/scholarpedia.7117 revision #135539 [link to/cite this article] Post-publication activity Curator: Carlo Rovelli Dr. Carlo Rovelli, Center of Theoretical Physics, and Marseille, France**

(2128) Quantum Gravity is the name given to any theory that describes gravity in (eb) the regimes where quantum effects cannot be disregarded.

(2129) At present, there is no such a theory which is universally accepted and confirmed by experience. Therefore the term "Quantum Gravity" indicates (eb) more an open problem than a specific theory.

(2130) Several research lines, at different levels of development, offer tentative solutions to the problem. These tentative quantum-gravity theories are variously viewed as competing research directions, or as contributions to the common goal of finding the physically correct theory. The quest for the good quantum theory of gravity bears on a number of fundamental issues and it is sometimes presented as the most important open problem in fundamental physics: the "Holy Grail" of contemporary theoretical physics.

**The problem of quantum gravity**

(2131) The current theory of (e) gravity is general relativity, a theory developed by Albert Einstein at the beginning of the XX century, which has changed in depth our understanding of (e) what are space and time.

(2132) Similarly, quantum mechanics has changed in depth our understanding of matter, energy and causality. These two theories, today very well empirically supported, have opened (eb) a major conceptual revolution in physics.

(2133) However, they appear to be incompatible, at least at a first reading, because (e) each of the two is formulated on the basis of (e) principles that are explicitly contradicted by (e) the other theory.

(2134) Therefore in spite of the immense development of scientific knowledge it has delivered, the physics of the XX century has left us with (e&eb) a major confusion on what are the basic conceptual ingredients for (e) understanding the physical world.

(2135) In other words, the scientific revolution opened by general relativity and quantum mechanics at the beginning of (e) the XX century is not concluded yet, and a new synthesis is required. Quantum gravity, merging general relativity and quantum mechanics should be this new synthesis.

**Quantum space and quantum time**

(2136) Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.

(2137) In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.

(2138) They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.

(2139) When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.

(2140) Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.

**The Planck scale**

(2141) Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".

(2142) Here $\hbar$ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.

(2143) The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.

(2144) Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)

(2145) Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.

## Physical applications

Some phenomena where quantum gravity plays a major role are the following.

**The microscopic structure of spacetime**

(2146) If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.

(2147) Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length

(2148) Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.

(2149) A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.

**Early cosmology**

(2150) According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.

(2151) However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.

(2152) A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.

(2153) Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".

(2154) One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.

**Black holes**

(2155) Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.

(2156) First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.

(2157) Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.

(2158) What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.

(2159) Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.

**Astrophysical effects**

(2160) Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.

(2161) An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.

(2162) The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.

**Main approaches**

(2163) The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.

(2164) Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.

(2165) The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.

**String theory**

(2166) String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.

(2167) Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.

(2168) String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.

(2169) Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.

(2170) However, when summed the entire series appears to be (=) divergent.

(2171) In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.

(2172) Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.

(2173) Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.

(2174) In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.

**Loop quantum gravity and spin foams**

(2175) Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.

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| (2176) The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity. |
| (2177) Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime. |
| (2178) The granular properties of space can be (=) explicitly computed. |
| (2179) In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory. |
| (2180) The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks. |
| (2181) Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure. |
| (2182) The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit. |
| (2183) Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism. |
| (2184) This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism. |

**Noncommutative geometry**

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| (2185) Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime. |
| (2186) Quantum mechanics teaches us that dynamical quantities are (=) noncommutative --meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed. |
| (2187) It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry. |
| (2188) A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces. |
| (2189) Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces. |
| (2190) This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity. |

**Others**

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| (2191) Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays |

(2192) . Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection.

(2193) Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible.

(2194) The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.

Let $p_8(t, a_8)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(g_8)Q_8$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_8(0) = 0, \quad g_8(P_8) > 0, \quad P_8 \text{ 0}, \qquad \lim_{P_8 \to \infty} g_8(P_8) = (K_8)_8 \text{ ∞}. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_8 + (\rho_a)_8 + D_8(a_8, P_8, Q_8)\rho_8 + 0, \qquad 0 < a_8 \text{ ∞},$$

$$\rho_8(t,0) = \int_0^\infty n_8(a_8, P_8, Q_8)\rho_8(t_8, a_8)da_8$$

Where $D_8$ is the mortality function and $n_8$ is the fertility function.

Put $_8$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_8}{dt}\right]_8 = x_8\left[\frac{dQ_8}{dt}\right]_8 - x_8\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_8}{dt}\right]_8 = x_3\left[\frac{dQ_3}{dt}\right]_8 - x_4\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_8}{dt}\right]_8 - x_6\left[\frac{dQ_3}{dt}\right]_8$$

Where $x_8, x_8, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_8$ and $g_8(P_8)Q_8$

Now when you add the term $(b_8Q_8)$ with $g_8(P_8)(Q_8)$ we are infect see is that $(b_8Q_8)$ is credited incrementally part of $g_8(P_8)(Q_8)$ or instant nearly, when subtracted from $\frac{dQ_8}{dt}$ , it means that $g_8(P_8)(Q_8)$ is slowly eaten by $\frac{dQ_8}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_8 = \mu_8(a_8) + (f_8)_8(P_8)_8 + (f_8)_8(Q)_8 \quad (2)$$

$$n_8 = \beta_8(a_8)h_8(Q_8), \quad (3)$$

Where $\mu_8$the age- dependent debit rate whit is $\mu_8(a_8)$ 0 $\forall a > 0$; $(f_8)_8$ is the density- dependent death rate with $f_8(0) = 0$, $(f_8^8)_8(P_8)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_8)_8(P_8) = (k_8)_8$ ; $f_8$ is the functional response of creditors to prey with $[f_8]_8(0) = 0, (f_8^8)_8(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_8)_8(Q_8) = [k_8]$ ; $\beta$ is the age

dependent birth rate with $\beta(a) \geq 0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^{8}(Q) \geq 0, \forall Q \geq 0$, and $\lim_{Q\to\infty} h(Q) = k_3 \geq 0$.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_8 + (\rho_a)_8 = -\mu_8(a_8) + (f_8)_8 P_8) + \big(f_8)_8(Q_8)\big)(\rho_8)_8 \quad (4)$$

$$\rho_8(t_8, 0) = h_8(Q_8) \int_0^{\infty} \beta_8 (a_8)\rho_8(t, a_8)da, \quad (5)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(P_8)Q_8 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_8$ and $Q_8$ of (4-6) satisfy

$$0 \leq \rho_8(t, a_8) \leq \rho_8^*(t, a_8),$$

$$0 \leq Q_8(t) \leq Q_8^*(t),$$

Where $\rho_8^*(t, a_8)$ and $Q_8^*(t)$, are solutions of

$$(\rho_t^*)_8 + (\rho_a^*)_8 = -(\mu_8(a_8) + (f_8)_8(P_8^*) + (f_8)_8 Q_8^*))\, \rho_8^*, \quad (7)$$

$$\rho_8^*(t, 0) = \int_0^{\infty} \beta_8(a_8)\, \rho_8^*(t, a_8)da_8 \quad (8)$$

$$\frac{dQ_8^*}{dt} = -b_8 Q_8^* + g_8(P_8^*)Q_8^*, \quad (9)$$

With $P_8^*(t) = \int_0^\infty \rho_8^*\ (t, a_8)da_8$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_8^*(t)}{dt} = P_8^*(t)\big(r_8 - (f_8)(P_8^*) - (f_8)_8(Q^*)\big), \tag{10}$$

$$\frac{dQ_8^*(t)}{dt} = Q_8^*(t)(-b_8 + g_8(P_8^*)), \tag{11}$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_8(a)e^{-\int_0^a \mu_8(\tau)dt}e^{-ra}\ da_8 = 1. \tag{12}$$

If $r < 0$, the trivial solution $P_8^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(t)dt}\ da_8 > 1, \tag{$H1$}$$

Which guarantees $r_8 > 0$.

*Remark.* If $\beta_8(a) = [\beta_0]_8 a_8 e^{-\alpha_8 a_8}, \mu_8(a_8) \equiv \mu_8$ then the solution of (12) is

$$r_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_8]_8(P_8) = 0, [f_8](Q_8) = \lambda_8 Q_8\ h_8(Q_8) = 1, \qquad and\ \ g_8(P_8) = c_8 P_8$$

$$[\rho_8]_8 + [\rho_a]_8 = -(\mu_8 + \lambda_8 Q_8)\rho_8$$

$$\rho_8(t,0) = \int_0^{\infty} [\beta_0]_8 a_8 e^{-\alpha_8 a_8} \rho_8(t,a_8)\, da_8$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + c_8 P_8 Q_8$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_8}{dt} = \omega_8 P_8 - \lambda_8 P_8 Q_8$$

$$\frac{dQ_8}{dt} = b_8 Q_8 - c_8 P_8 Q_8$$

Where $\omega_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$

## 35. THE EQUILIBRIUM $(\rho_0)_8(a_8), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_8(a_8), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_8 + (\rho_a)_8 = -\big(\mu_8(a_8) + (f_8)_8(P_8)\big)\rho,$$

$$\rho_8(t,0) = \int_0^{\infty} \beta_8(a_8)\rho_8(t,a_8)da_8 \qquad (13)$$

The solution of (13) has the form of

$$\rho_8(t,a_8) = Ce_8^{r(t-a_8)} - \int_0^a (\tau)d\tau - \int_0^8 (f_8)_8\big(P_8(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_8(t) = C_8 e^{r_8 t - \int_0^t (f_8)_8\left(P_8(\tau)\right)d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau}\, da.$$

Hence, there exists an equilibrium $P_0 = \int_8^{-8}(r)$ if $k_8 > r$. This is equivalent to

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{(k_8)_8 a_8} da_8 < 1.$$

In fact, if we define

$$R_8(s_8) = \int_0^\infty \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau)d\tau} e^{-s_8 a_8}\, da_8$$

$R(s)$ is a decreasing function. $R_8(r_8) = 1$ And $R_8(k_8)_8 < 1$ imply $r_8 < (k_8)_8$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_8(a_8), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{k_8 a_8} da_8 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_8 \equiv \int_0^\infty (\rho_0)_{88}(a_8)da_8$ *satisfies*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{-(f_8)_8(P_0)a_8} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_8(a_8), 0)$

Let $u_8(t,a_8) = \rho(t,a_8) - \rho_0(a_8)$ Then

$$N_8(t) \equiv \int_0^\infty u_8(t,a_8)da_8 = \int_0^\infty \rho_8(t,a_8)da_8 - \int_0^\infty (\rho_0)_8(a_8)da = P_8(t) - P_0)_8$$

And

$$(u_8)_8 + (u_a)_8 = (\rho_t)_8 + (\rho_a)_8 - (\rho_0^8)_8(a_8)$$

$$= -\mu_8(a_8) + (f_8)_8(P_8) + (f_8)_8(Q_8))\,(u_8 + (\rho_0)_8) + \mu_8(a_8)(\rho_0)_8 + (f_8)_8(P_0)_8(\rho_0)_8$$

$$= -\mu_8(a_8)u_8 - (f_8)_8(P_8) + (f_8)_8(Q_8))\,u_8 - (\rho_0)_8\,(f_8)_8(P_8) - (f_8)_8(P_0)_8$$

$$+ (f_8)_8(Q_8)), \qquad (15)$$

Where $(\rho_0^8)_8(a_8)$ denotes $d(\rho_0)_8(a_8)/da_8$

Expansion $(f_8)_8(P_8)$ and $(f_8)_8(Q_8)$ at $(P_0)_8$ and 0 yield respectively

$$(f_8)_8(P_8) = (f_8)_8((P_0)_8 + N_8)$$

$$= (f_8)_8(P_0)_8 + (f_8^8)_8(P_0)_8N_8 + o_8(N_8), \qquad (16)$$

$$(f_8)_8(Q_8)$$

$$= (f_8^8)_8(0)Q + o_8(Q_8) \qquad (17)$$

Write $f_8(P_0) = f_{80}, f_8'(P_0) = f_{80}'$, and $f_8'(0) = f_{80}'$ and substitute (16) and (17) into (15) to obtain

$$(u_8)_8 + (u_a)_8 \approx -(\mu_8(a_8) + (f_{80})_8)u - (f_{80}^8)_8(\rho_0)_8(N_8) - (f_{80}^8)_8(\rho_0)_8(Q_8)$$

$$u_8(t,0) = \rho_8(t,0)(\rho_0)_8(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a) + (\rho_0)_8(a_8)da_8 - (\rho_0)_8\,(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a_8)da_8 + (h_8(Q_8) - 1)(\rho_0)_8\,(0) \qquad (18)$$

Expanding

$$h_8(Q_8) = 1 + (h^8)_8(0)Q_8 + 0(Q_8) \equiv 1 + h_0^{'}Q_8 + 0Q_8$$

And substituting it into (18), we arrive at

$$u_8(t,0) \approx \int_0^\infty \beta_8\,(a_8)u_8(t,a_8)da_8 + (h_{a_8}^{'})(\rho_0)_8 Q_8$$

The linearization is the following :

$$[u_8] + [u_8] = -(\mu_8(a_8) + [t_{80}]_8 u_8 - [\alpha_8]_8(a_8)N_8 - [\alpha_8]_8(a_8)Q_8$$

$$u_8(t,0) = \int_0^\infty \beta_8(a_8)u_8(t,a_8)da_8 + [\xi_8]Q_8 \qquad (19)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8$$

where

$$[\alpha_8]_8(a_8) = [t_8^8 a]_8[p,o]_8(a_8)d2(a_8) = [f_{80}^8]_8(\rho\, o\,)_8(a_8)\xi = [\mathrm{h\ o}^8]_8[\rho \mathrm{o}]_8(\mathrm{o}) \text{ and } \mathrm{b}_8 = \mathrm{b}_8 - g_8(P_8)$$

If $g_8(P_0) > b$, clearly, $(\rho_0)_8(a_8), 0)$ is unstable and the predator population is persistent. Assume $g_8(P_0)_8 < b$.

Then, $Q_8(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_8 + u_a = -(\mu_8(a_8) + (f_{80})_8)u_8 - (\alpha_8)_8(a_8)N_8$$

$$= -(\mu_8(a_8) + f_{80})u_8 - \alpha_8(a_8)\int_0^\infty u_8(t,a_8)da_8$$

$$u_8(t,0) = \int_0^\infty \beta_8\,(a_8)u_8(t,a_8)da_8$$

Using the method developed by Cushing [2, 3], we suppose

$$u_8(t,a_8) = \tilde{u}_8(a_8)^{e^{\gamma(t-a_8)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_8}\,da_8$. Then,

A solution of the first differential equation for $\tilde{u}(a_8)$ gives,

$$\tilde{u}_8(a_8) = e^{-\int_0^a \mu_8(\tau_8)d\tau - (f_{80})_8 a}\left(\tilde{u}_8(0) - w\int_0^a e^{\int_0^\tau \mu_8(s_8)ds + (f_{80})}\tau_8\,(\alpha_8)_8(\tau_8)e^{\gamma_8\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\right)\tilde{u}_8(0)$$

$$+w\int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\int_0^a (\alpha_8)_8(\tau_8)e^{-\gamma_8\tau_8}e^{(f_{80})_8\tau_8 + \int_0^\tau \mu_8(s_8)ds}d\tau_8\,da_8$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}da\right)(\tilde{u}_8(0))_8$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_8 a_8-(f_{80})_8 a_8-\int_0^a \mu_8(\tau_8)d\tau}\int_0^a e^{(f_{80})_8\tau_8+\int_0^\tau \mu_8(s_8)ds+y\tau}\alpha_8(\tau)d\tau\, da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_8 a_8-(f_{80})_8 a_8-\int_0^a \mu_8(t_8)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_8\,(a_8)Y_8(\gamma_8,a_8)da_8\right)u_8(0)$$

$$+w\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)\int_0^a(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8=0,$$

(24)

$$\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\tilde{u}_8(0)$$

$$+\left(-1-\int_0^{\infty}Y_8(\gamma_8,a_8)\int_0^a(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8\right)w=0.$$

Equation (20) has a solution of the form of $u_8(t,a_8)=u_8(a_8)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_8(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_8(\gamma_8,a_8)=\int_0^a(\rho_0)_8(\tau_8)Y^{-8}\,(\gamma_8,a_8)d\tau_8$. We obtain the characteristic equation

$$f_{80_8}^{l}\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\left(\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$+\left(1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)da\right)\left(1+(f_{80}^{l})_8\int_0^{\infty}Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_8(a_8) = (\rho_0)_8(0)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - (f_{80})_8 a_8}$$

$$I_8(\gamma_8, a_8) = \int_0^a (\rho_0)_8(\tau_8)e^{\gamma_8\tau_8 + (f_{80})_8\tau + \int_0^\tau \mu_8(s_8)ds_8}d\tau$$

$$= \frac{(\rho_0)_8(0)}{\gamma_8}(e^{\gamma_8 a_8} - 1)$$

Let

$$E_8(a_8) = e^{-(f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau_8}$$

Then $[Y_8]_8[\gamma_8]_8, a_8) = E_8(a_8)e^{-\gamma_8 a_8}$. by a simple calculation,

$$\int_0^\infty \beta_8(a_8)Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma}\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}\ da\right);$$

$$\int_0^\infty Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma_8}\left(\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}\ da_8\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_8(0)}{\gamma_8}[f_{80}^8]\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}da_8\ (\int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8)$$

$$+\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right).$$

$$\left(1+\frac{(\rho_0)_8(0)}{\gamma_8}(f_{80}^{'})_8\left(\int_0^{\infty}E_8(a_8)da_8-\int_0^{\infty}E_8(a_8)e^{-\gamma_8 a_8}da_8\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right)(1+\frac{(f_{80}^{'})_8}{\gamma_8}(\rho_0)_8(0)\int_0^{\infty}E_8(a_8)da_8)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)da_8=1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\neq 0,$$

Since $\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)\int_0^{a}\alpha_8(\tau_8)Y^{-8}(\gamma_8,\tau_8)d\tau_8 da_8>0$. Therefore, the characteristic solution is

$$\gamma_8=-(f_{80}^{'})_8(\rho_0)_8\int_0^{\infty}E_8(a_8)da_8.$$

Since $f_8^{'}(P_0)_8>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_8),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_8),0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_8 + (\mu_8(a_8) + [f_8]_8(P_8^*) + [f_8]_8(Q_8^*))P_8^* = 0, \quad (26)$$

$$p_8^*(0) = h_8(Q_8^*)\int_0^\infty \beta_8(a_8)p_8^*(a_J)p_8^*(a_8)da_8 \quad (27)$$

$$g_8(P_8^*) - b_8 = 0, \quad (28)$$

Where $P_8^* = \int_0^\infty p_8^*(a_8)\, da_8$

Equation (28) yields $P^* = g^{-8}(b)$,

From (26),

$$p_8^*(a) = p_8^*(0)e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}$$

Substituting into (27), we need to solve

$$\left(1 - h_8(Q_8^*)\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}da_8\right)p_8^*(0) = 0$$

For $Q_8^*$

Fix $p_8^*(0)$ and $P_8^*$ Let

$$H_8(z_8) = 1 - h_8(z_8)\int_0^\infty \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}da_{,8}$$

And assume that

$$[k_3]_8 \int_0^{\infty} \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(g_8^{-8}(b_8))a_8 - [k_8]_8 a_8}\, da_8 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty} [H_8(z_8)] > 0.$$

Clearly,

$$\frac{dH_8}{dz_8} = h_8(z)(\,[f_8^{'}]_8(z_8)\,) \int_0^{\infty} a_8\beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - f_8(P_8^*)a_8 - f_8(z_8)a_8} da_8$$

$-h^{'}(z_8)\int_0^{\infty} \beta_8(a_8)\, e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}\ da_8$

By the assumptions on $\left[f_8^8\right]_8 and\ [h^8]^8 dH_8/dz > 0$. Hence, if

$$H_8(0) = 1 - \int_0^{\infty} \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8} da_8 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_8(0) < 0$ follows if $g_8^-(b_8) < [f_8^{-8}]_8(r_8)$, where $r_8$ is the solution of (12).

Indeed, since

$$R_8(s_8) = \int_0^{\infty} \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 e^{-y_8^a}} da_8$$

Is a decreasing function of $s_8$ when $P_8^* = [g^{-8}]_8(b_8) < [f_8^{-8}]_8(r_8)$, that is $[f_8]_8(P_8^*) < r$, $R\big(f_8(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_8(g^{-8}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8\left(g_8^{-8}(b)\right)a_8} > 1. \qquad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_8^*(a_8), Q_8^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_8^*(a_8), Q_8^*)$ is more complicated. We again consider perturbations $m_8(t_8 a_8) = p_8(t_8, a_8) - p_8^*(a_8)$ and $W_8(t) = Q_8(t) - Q^*$.Then

$M_8(t) \equiv \int_0^\infty m_8(t, a_8)da_8 = P_8(t) - P_8^*,$

$$[m_8]_8 + [m_a]_8 = [p_8]_8 + [p_a]_8 - \left[p^*\right]_8(a_8) \approx -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8$$

$$-[f_{8*}^8]_8 p_8^*(a_8)M - [f_{8*}^8]_8 p_8^*(\mathrm{a})W_8,$$

$$m_8(t,0)p_8(t,0) - p_8^*(0) \approx h_* \int_0^\infty \beta_8(a_8)m_8(t,a_8)\,da + \frac{[h_*^8]_8}{[h_*]_8}\, p_8^*(0)W_8$$

$$\frac{dW_8}{dt} = \frac{dQ_8}{dt} \approx [g_*^8]_8 Q_8^* M_{,8}$$

Where $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $\left[f_{8*}^{'}\right]_8$ denoting $\left[f^{'}{}_8(P^*)\right]_8\left[f_{8*}^{'}\right]_8$ denoting $[f_8]_8\,(Q_8^*,\quad h_{*}^{'}$ denoting $h_*^{'}$ denoting $h_8^{'}\,(Q_8^*)$, and $\left[g_*^{'}\right]_8$

Denoting $\left[\dot{g}\right]_8(P_8^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_8]_8 + [m_a]_8 = -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8 - \left[\dot{f}_{8*}\right]_8[p_8^*](a_8)M_8 - \left[\dot{f_{8*}}\right]_8 p_8^*(a_8)W_8$$

$$m_8(t,0) = [h_*]_8 \int_0^\infty \beta_8(a_8)m_8(t,a_8)da_8 + \frac{[h_*^8]_8}{[h^*]_8}[p^*]_8(0)\, W_{,8} \qquad (30)$$

$$\frac{dW_8}{dt} = [g_*^8]_8[Q_8^*]M_8$$

Integrating $[m_8]_8(t,a_8)$ with respect to a from 0 to ∞, we have

$$\frac{dM_8(t)}{dt} = \int_0^\infty [mt]_8(t,a_8)da_8$$

$$= -\int_0^\infty [m_a]_8\,(t,a_8)da_8 - \int_0^\infty \mu(a_8)m(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + \left[\dot{f}_{8*}\right]_8 \int_0^\infty p_8^*(a_8)da_8\right)M_8 - [f_{8*}^8]_8 \int_0^\infty p_8^*(a_8)da_8 W_8$$

$$= m_8(t,0) - \int_0^\infty \mu_8(a_8)\, m_8(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + \left[\dot{f_{8*}}\right]_8 P_8^*\right)M_8 - \left[\dot{f_{8*}}\right]_8 P_8^* W_{,8}$$

From (30), we arrive at the equations

$$\frac{dW_8(t)}{dt} = \int_0^\infty ([h_*]_8\beta_8(a_8) - \mu_8(a_8))m_8(t,a_8)da_8$$

$$-\left([f_{8*}] + [f_{8*}]_8 + [f_{8*}']_8 P_8^*\right) M_8 + \left(\frac{[h_*']_8}{[h_*]^8}\, P_8^*(0) - f_{8*}' P_8^*\right) W_8,$$

(31)

$$\frac{dW_8(t)}{dt}\left[g'_*\right]_8 Q_8^*\, M_{.8}$$

By the mean value theorem,

$$\int_0^\infty ([h_{"*}]_8 \beta_8(a_8) - \mu_8(a_8)) m_8(t, a_8) da_8 = ([h_*]_8 \beta_8(\theta_8) - \mu_8(\theta_8)) M_{,8}$$

Where $0< \theta_8 = \theta_8(t) < \infty$.

And if a $^8_8$[a $^8_8$]2<4 [a $_8$]2<4 a $_8$e,

Re $\lambda_8[\lambda_8]_8 = \text{Re}\lambda_8[\lambda_8]_8 = -\frac{8}{8}a_8[\text{a2}\ ]_8 < 0$.

Let ⊿=max$_{8<u<\infty}$ |h $_*$ [h $_*$] $_8$β $_8$($a_8$) − $\mu_8(a_8)$|. Then|$n_8$(t)| ⊿∀t . Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium(ρ $^*$ (a), $\boldsymbol{Q}^*$) is asymptotically stable if either a $^8_8$ <4 a $_8$ e and [a $_8$] > 1⊿,or [a $^8_8$] ⊿[a $_8$]e and (a $_8$ − $\sqrt{(a^8_8)_8 - 4(a_8)_8 e}$ > 1⊿

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+8}\right)$ and β $_8$($a_8$) = ae-$^{ma}$ Then ⊿=[$\mu_0$] + 2 Simple calculation leads to

$$\int_0^{\infty} \beta_8(a_8)\mathrm{e}^{-\int_0^{\mathrm{a}} \mu_8(\tau_8)+d\tau_8}\mathrm{da} = \frac{1}{(\mathrm{m}+\mu_0+1)^8}+\frac{1}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_8 + [\mu_0]$ 0.5. Assume$[f_8](x_8)= K_8x_8$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_8(x_8)= [0x_8]/(2 + x_8)$; the functional responses $[[f_8](x) = [0]\, x_8/(1 + x_8)$and $h_8(x_8)$=$(1+\omega_8x_8)/(1 + x_8)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_8 + [\mu_0] + [\mathrm{K}][P_0])}{[m_8]_8 + (\mu_0)_8 + [\mathrm{K}][P_0])_8^3\ [\mathrm{K}]_8[P_0]_8^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_8$i e $[P_0] < b/(\delta_8 - b_8)$then the equilibrium$(P_0)$ $[P_0]_8(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_8$

$[f_{8*}] = K_8\ \mathrm{P}^*\ [f_{8*}] = \frac{\sigma \mathrm{Q}_8^*}{8+\mathrm{Q}_8^*}$ $[h_*] = \frac{8+\omega_8\mathrm{Q}_8^*}{8+\mathrm{Q}_8^*}$

$[f_{8*}] = K_8\ [\mathrm{f}_{8*}^{8}] = \frac{\sigma_8}{(8+\mathrm{Q}_8^*)^8}[\ \mathrm{g}_*^{8}] = \frac{\sigma_8}{(8+\mathrm{P}_8^*)^8}$ $[\mathrm{h}_*^{8}] = \frac{\omega_8-8}{(8+\mathrm{Q}_8^*)^8}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_8}{(m + [\mu_0]_8 + 1 + \frac{K_8b_8}{\delta_8 + b_8} + \sigma)^8} + \frac{1\omega_8}{([m_8]_8 + [\mu_0]_8 + 1 + \frac{K_8b_8}{\delta_8 + b_8} + \sigma_8)^3} < 1$$

And

$$\frac{1}{(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^8} + \frac{1\omega_8}{(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8} + (m_8 + [\mu_0]_8 + \frac{K_8 b_8}{\delta_8 - b_8} + \sigma_8)\, Q_8^*)^3$

$= m_8 + [\mu_0]_8 + 3 + \frac{K_8 b_8}{\delta_8 - b_8}(1 + Q^*)(1 + \omega_8 Q_8^*$

Now

$$\left[a_{11}\right] - 4\left[a_{88}\right] e = \left(\frac{1K_8 b_8}{\delta_8 - b_8}\right)^8 + \frac{4K_8 b_8 \sigma_8 Q_8^*}{(\delta_8 - b_8)(8 + Q_8^*)} + \frac{\sigma_8^8 Q_8^{*8}}{(1 + Q_8^*)^8}$$

$$\frac{4\, b_8 \sigma_8 (\delta_8 - b) Q^*}{\delta_8 (1 + Q_8^*)^8} - \frac{4\, (\delta_8 - b_8)^8 (1 - \omega_8) Q_8^*}{\delta_8 (1 + \omega_8 Q_8^*)(1 + Q_8^*)} P^*(0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_8\, ae_8^{mu}\ [\mu_0]_8 - \frac{a_8}{a_8 + 8}$$

Hence if $[a_8^8 > 4(a_8)_8 e$ and

$$[a_8] = \frac{1K_8 b_8}{\delta_8 - b_8} + \frac{\sigma_8 Q_8^*}{1 + Q_8^*} > 1\Delta$$

**The nontrivial equilibrium is stable.**

IF we define

$$[a_8]_8 = [f_{8*}]_8 + [f_{8*}]_8 + [f_{8*}] P_8^*$$

$$[a_8]_8 = \left[f_{8*}^{8}\right] P_8^* - \frac{h_*^8}{h_*} P^*(0)$$

$$e = [g_*^8]_8 Q_8^*$$

$$\eta_8(t) = \left[h_*\right]_8 \beta_8 \left(\theta_8(t)\right) - \mu_8(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\left[a_8\right] M_8 - [a_8]_8 W_8 - [n(t)] M_8$$

$$w_8(t) = eM_8$$

Write

$$A = \begin{pmatrix} -[a_8]_8 - [a_8]_8 \\ e \qquad 0 \end{pmatrix}$$

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| $B=(t)=\begin{pmatrix} [\eta(t)]_8 & 0 \\ 0 & 0 \end{pmatrix}$ |
| And |
| $X(t)=\begin{pmatrix} M_8 \\ W_8 \end{pmatrix}$ |
| The solution of (32)is |
| $X_8(t) = e^{A,t}[X]_8(0) + \int_0 e^{At}$ |
| This equation notwithstanding its importance in all the models stand out for its importance for the following models<br><br>to which it is applied: |
| **Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016** |

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| (270)These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field. |
| (271)The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets. |
| (272) Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state. |
| (273) Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands. |
| (274)The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI |
| (275) The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9]. |
| (276)This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive. |

(277)More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

(278)Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.

(279)Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.

(280) The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.

(281)This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.

(282)The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.

(283)This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.

(284) For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.

(285)The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(286)Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.

(287)We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.

(288)When this is true—as is often the case—the system serves as (=) a "bath" for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called **thermalization.**

(289)However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.

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| (290)Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1). |
| (291)Bringing them to thermal equilibrium requires (e) an external bath. |
| (292) But so far, there has been no rigorous proof of (e) the existence of **such many-body localization**. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption. |
| (293)The concept of **localization** was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations). |
| (294)The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century. |
| (295)But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information. |
| (296)Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment. |
| (297)For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood. |
| (298) It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7]. |
| (299) Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature. |
| (300) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators. |
| (301)However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8]. |
| (302)Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems. |
| (303) Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL. |
| (304)As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked. |
| (305)His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice. |
| (306) An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site. |

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| (307)This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself). |
| (308)This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition. |
| (309)However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime: |
| (310) These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects. |
| (311)It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies). |
| (312) This is a realistic scenario for (e) all known systems. |
| (313)On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects. |
| (314)On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other. |
| (315)No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true. |
| (316) His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found. |
| (317)Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization? |
| (318)Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects? |
| (319)Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions. |
| (320)So in those systems, the quantum phase transition does indeed exist. |
| (321)But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension. |
| (322)Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters. |

# PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION ATTRITION MODELS

## MODULE NINE

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Model is applied to following formulations Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society**

**Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice ($\alpha$=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP This paper is divided into four parts. In part A, some general considerations about (e&eb) gravitational radiation are followed by (e) a treatment of the scalar wave equation in the manner later to be applied to (e&eb) Einstein's field equations. In part B, a co-ordinate system is specified which is suitable for (e) investigation of outgoing gravitational waves from (e) an isolated axi-symmetric reflexion-symmetric system. The metric is expanded in (eb) negative powers of (e) a suitably defined radial co-ordinate r, and the vacuum field equations are investigated in detail. It is shown that the flow of information to infinity is controlled by (e&eb) a single function of (e) two variables called the news function. Together with initial conditions specified on a light cone, this function fully defines (eb) the behaviour of the system. No constraints of (e) any kind are encountered. In part C, the transformations leaving (eb) the metric in the chosen form are determined. An investigation of the corresponding transformations in (eb) Minkowski space suggests (eb) that no generality is lost by (e) assuming that the transformations, like the metric, may be expanded in (eb) negative powers of r. Gravitational Waves in General Relativity VII Waves from Axi-Symmetric Isolated Systems Royal Society Proceedings H. Bondi, M. G. J. van der Burg, A. W. K. Metzner Published 21 August 1962.DOI: 10.1098/rspa.1962.0161 STOP There is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant $\Lambda$, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive $\Lambda$ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the $\Lambda>0$ generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible. Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of (e) GW signals from BBHs with component spins $\chi_{1,2}\in[-0.99,0.99]$ aligned (e&eb) with the orbital angular momentum, component masses $m_{1,2}\in[2,48]M_\odot$, and (e&eb) total mass $M_{total}\leq 50\ M_\odot$. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin**

**bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins >+0.6 and component masses ≳20 M⊙. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201**

## 36. MODEL FORMULATION

Let $p_9(t, a_9)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(g_9)Q_9$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_9(0) = 0, \quad g_9(P_9) > 0, \ P_9 \ 0, \qquad \lim_{P_9 \to \infty} g_9(P_9) = (K_9)_9 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_9 + (\rho_a)_9 + D_9(a_9, P_9, Q_9)\rho_9 + 0, \qquad 0 < a_9 \ \infty,$$

$$\rho_9(t,0) = \int_0^\infty n_9(a_9, P_9, Q_9)\rho_9(t_9, a_9)da_9$$

Where $D_9$ is the mortality function and $n_9$ is the fertility function.

Put $_9$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_9}{dt}\right]_9 = x_9\left[\frac{dQ_9}{dt}\right]_9 - x_9\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_9}{dt}\right]_9 = x_3\left[\frac{dQ_3}{dt}\right]_9 - x_4\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_9}{dt}\right]_9 - x_6\left[\frac{dQ_3}{dt}\right]_9$$

Where $x_9, x_9, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_9$ and $g_9(P_9)Q_9$

Now when you add the term $(b_9Q_9)$ with $g_9(P_9)(Q_9)$ we are infect see is that $(b_9Q_9)$ is credited incrementally part of $g_9(P_9)(Q_9)$ or instant nearly, when subtracted from $\frac{dQ_9}{dt}$ , it means that $g_9(P_9)(Q_9)$ is slowly eaten by $\frac{dQ_9}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_9 = \mu_9(a_9) + (f_9)_9(P_9)_9 + (f_9)_9(Q)_9 \quad (2)$$

$$n_9 = \beta_9(a_9)h_9(Q_9), \quad (3)$$

Where $\mu_9$the age- dependent debit rate whit is $\mu_9(a_9)$ 0 $\forall a > 0$; $(f_9)_9$ is the density- dependent death rate with $f_9(0) = 0$, $(f_9^9)_9(P_9)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_9)_9(P_9) = (k_9)_9$ ; $f_9$ is the functional response of creditors to prey with $[f_9]_9(0) = 0, (f_9^9)_9(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_9)_9(Q_9) = [k_9]$ ; $\beta$ is the age dependent birth

rate with $\beta(a) \geq 0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h'(Q) \leq 0, \forall Q \geq 0$, and $\lim_{Q\to\infty} h(Q) = k_3 \geq 0$.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_9 + (\rho_a)_9 = -\mu_9(a_9) + (f_9)_9 P_9) + \big(f_9)_9(Q_9)\big)(\rho_9)_9 \tag{4}$$

$$\rho_9(t_9, 0) = h_9(Q_9) \int_0^\infty \beta_9(a_9)\rho_9(t, a_9)da, \tag{5}$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(P_9)Q_9 \tag{6}$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_9$ and $Q_9$ of (4-6) satisfy

$$0 \leq \rho_9(t, a_9) \leq \rho_9^*(t, a_9),$$

$$0 \leq Q_9(t) \leq Q_9^*(t),$$

Where $\rho_9^*(t, a_9)$ and $Q_9^*(t)$, are solutions of

$$(\rho_t^*)_9 + (\rho_a^*)_9 = -(\mu_9(a_9) + (f_9)_9(P_9^*) + (f_9)_9 Q_9^*))\, \rho_9^*, \tag{7}$$

$$\rho_9^*(t, 0) = \int_0^\infty \beta_9(a_9)\, \rho_9^*(t, a_9)da_9 \tag{8}$$

$$\frac{dQ_9^*}{dt} = -b_9 Q_9^* + g_9(P_9^*)Q_9^*, \tag{9}$$

With $P_9^*(t) = \int_0^\infty \rho_9^*\ (t, a_9)da_9$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_9^*(t)}{dt} = P_9^*(t)\big(r_9 - (f_9)(P_9^*) - (f_9)_9(Q^*)\big), \quad (10)$$

$$\frac{dQ_9^*(t)}{dt} = Q_9^*(t)(-b_9 + g_9(P_9^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_9(a)e^{-\int_0^a \mu_9(\tau)dt}e^{-ra}\ da_9 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_9^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(t)dt}\ da_9 > 1, \quad (H1)$$

Which guarantees $r_9 > 0$.

*Remark.* If $\beta_9(a) = [\beta_0]_9 a_9 e^{-\alpha_9 a_9}, \mu_9(a_9) \equiv \mu_9$ then the solution of (12) is

$$r_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_9]_9(P_9) = 0, [f_9](Q_9) = \lambda_9 Q_9\ h_9(Q_9) = 1, \quad and\ \ g_9(P_9) = c_9 P_9$$

$$[\rho_9]_9 + [\rho_a]_9 = -(\mu_9 + \lambda_9 Q_9)\rho_9$$

$$\rho_9(t,0) = \int_0^{\infty} [\beta_0]_9 a_9 e^{-\alpha_9 a_9} \rho_9(t,a_9)\, da_9$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + c_9 P_9 Q_9$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_9}{dt} = \omega_9 P_9 - \lambda_9 P_9 Q_9$$

$$\frac{dQ_9}{dt} = b_9 Q_9 - c_9 P_9 Q_9$$

Where $\omega_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$

THE EQUILIBRIUM $(\rho_0)_9(a_9), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_9(a_9), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_9 + (\rho_a)_9 = -\big(\mu_9(a_9) + (f_9)_9(P_9)\big)\rho,$$

$$\rho_9(t,0) = \int_0^{\infty} \beta_9\,(a_9)\rho_9(t,a_9)da_9 \qquad (13)$$

The solution of (13) has the form of

$$\rho_9(t,a_9) = Ce_9^{r(t-a_9)} - \int_0^a (\tau)d\tau - \int_0^9 (f_9)_9\big(P_9(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_9(t) = C_9 e^{r_9 t - \int_0^t (f_9)_9(P_9(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_9^{-9}(r)$ if $k_9 > r$. This is equivalent to

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{(k_9)_9 a_9} da_9 < 1.$$

In fact, if we define

$$R_9(s_9) = \int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau)d\tau} e^{-s_9 a_9}\, da_9$$

$R(s)$ is a decreasing function. $R_9(r_9) = 1$ And $R_9(k_9)_9 < 1$ imply $r_9 < (k_9)_9$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_9(a_9), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{k_9 a_9} da_9 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_9 \equiv \int_0^\infty (\rho_0)_{99}(a_9)da_9$ *satisfies*

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{-(f_9)_9(P_0)a_9} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_9(a_9), 0)$

Let $u_9(t,a_9) = \rho(t,a_9) - \rho_0(a_9)$ Then

$$N_9(t) \equiv \int_0^\infty u_9\,(t, a_9)da_9 = \int_0^\infty \rho_9(t, a_9)da_9 - \int_0^\infty (\rho_0)_9(a_9)da = P_9(t) - P_0)_9$$

And

$$(u_9)_9 + (u_a)_9 = (\rho_t)_9 + (\rho_a)_9 - (\rho_0^9)_9(a_9)$$

$$= -\mu_9(a_9) + (f_9)_9(P_9) + (f_9)_9(Q_9))\,(u_9 + (\rho_0)_9) + \mu_9(a_9)(\rho_0)_9 + (f_9)_9(P_0)_9(\rho_0)_9$$

$$= -\mu_9(a_9)u_9 - (f_9)_9(P_9) + (f_9)_9(Q_9))\,u_9 - (\rho_0)_9\,(f_9)_9(P_9) - (f_9)_9(P_0)_9$$

$$+ (f_9)_9(Q_9)), \qquad (15)$$

Where $(\rho_0^9)_9(a_9)$ denotes $d(\rho_0)_9(a_9)/da_9$

Expansion $(f_9)_9(P_9)$ and $(f_9)_9(Q_9)$ at $(P_0)_9$ and 0 yield respectively

$$(f_9)_9(P_9) = (f_9)_9((P_0)_9 + N_9)$$

$$= (f_9)_9(P_0)_9 + (f_9^9)_9(P_0)_9N_9 + o_9(N_9), \qquad (16)$$

$$(f_9)_9(Q_9)$$

$$= (f_9^9)_9(0)Q + o_9(Q_9) \qquad (17)$$

Write $f_9(P_0) = f_{90}, f_9'(P_0) = f_{90}'$, and $f_9'(0) = f_{90}'$ and substitute (16) and (17) into (15) to obtain

$$(u_9)_9 + (u_a)_9 \approx -(\mu_9(a_9) + (f_{90})_9)u - (f_{90}^9)_9(\rho_0)_9(N_9) - (f_{90}^9)_9(\rho_0)_9(Q_9)$$

$$u_9(t, 0) = \rho_9(t, 0)(\rho_0)_9(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t, a) + (\rho_0)_9(a_9)da_9 - (\rho_0)_9\,(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t, a_9)da_9 + (h_9(Q_9) - 1)(\rho_0)_9\,(0) \qquad (18)$$

Expanding

$$h_9(Q_9) = 1 + (h^9)_9(0)Q_9 + 0(Q_9) \equiv 1 + h_0'Q_9 + 0Q_9$$

And substituting it into (18), we arrive at

$$u_9(t,0) \approx \int_0^\infty \beta_9(a_9)u_9(t,a_9)da_9 + (h'_{a_9})(\rho_0)_9Q_9$$



PREDATOR-PREY POPULATION MODELS

The linearization is the following :

$$[u_9] + [u_9] = -(\mu_9(a_9) + [t_{90}]_9u_9 - [\alpha_9]_9(a_9)N_9 - [\alpha_9]_9(a_9)Q_9$$

$$u_9(t,0) = \int_0^\infty \beta_9(a_9)u_9(t,a_9)da_9 + [\xi_9]Q_9 \qquad (19)$$

$$\frac{dQ_9}{dt} = -b_9Q_9$$

where

$$[\alpha_9]_9(a_9) = [t_9^9a]_9[p,o]_9(a_9)d2(a_9) = [f_{90}^9]_9(\rho\, o\,)_9(a_9)\xi = [\mathrm{h\ o}^9]_9[\rho\mathrm{o}]_9(\mathrm{o}) \text{ and } \mathrm{b}_9 = \mathrm{b}_9 - g_9(P_9)$$

If $g_9(P_0) > b$, clearly, $(\rho_0)_9(a_9), 0)$ is unstable and the predator population is persistent. Assume $g_9(P_0)_9 < b$.

Then, $Q_9(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_9 + u_a = -(\mu_9(a_9) + (f_{90})_9)u_9 - (\alpha_9)_9(a_9)N_9$$

$$= -(\mu_9(a_9) + f_{90})u_9 - \alpha_9(a_9)\int_0^\infty u_9(t,a_9)da_9$$

$$u_9(t,0) = \int_0^\infty \beta_9\,(a_9)u_9(t,a_9)da_9$$

Using the method developed by Cushing [2, 3], we suppose

$$u_9(t,a_9) = \tilde{u}_9(a_9)^{e^{\gamma(t-a_9)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_9}\,da_9$. Then,

A solution of the first differential equation for $\tilde{u}(a_9)$ gives,

$$\tilde{u}_9(a_9) = e^{-\int_0^a \mu_9(\tau_9)d\tau-(f_{90})_9 a}\left(\tilde{u}_9(0) - w\int_0^a e^{\int_0^\tau \mu_9(s_9)ds+(f_{90})}\tau_9\,(\alpha_9)_9(\tau_9)e^{\gamma_9\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\right)\tilde{u}_9(0)$$

$$+w\int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\int_0^a (\alpha_9)_9(\tau_9)e^{-\gamma_9\tau_9}e^{(f_{90})_9\tau_9+\int_0^\tau \mu_9(s_9)ds}d\tau_9\,da_9$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}da\right)(\tilde{u}_9(0))_9$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\int_0^a e^{(f_{90})_9\tau_9+\int_0^\tau \mu_9(s_9)ds+y\tau}\alpha_9(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(t_9)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_9\,(a_9)Y_9(\gamma_9,a_9)da_9\right)u_9(0)$$

$$+w\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)\int_0^a(\alpha_9)_9(\tau)_9(Y^{-9})_9(\gamma_9,\tau_9)d\tau_9\,da_9=0,$$

(24)

$$\left(\int_0^{\infty}Y_9(\gamma_9,a_9)da_9\right)\tilde{u}_9(0)$$

$$+\left(-1-\int_0^{\infty}Y_9(\gamma_9,a_9)\int_0^a(\alpha_9)_9(\tau)_9(Y^{-9})_9(\gamma_9,\tau_9)d\tau_9\,da_9\right)w=0.$$

Equation (20) has a solution of the form of $u_9(t,a_9)=u_9(a_9)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_9(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_9(\gamma_9,a_9)=\int_0^a(\rho_0)_9(\tau_9)Y^{-9}\,(\gamma_9,a_9)d\tau_9$. We obtain the characteristic equation

$$f_{90_9}^{l}\left(\int_0^{\infty}Y_9(\gamma_9,a_9)da_9\right)\left(\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)I_9(\gamma_9,a_9)da_9\right)$$

$$+\left(1-\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)da\right)\left(1+(f_{90}^{l})_9\int_0^{\infty}Y_9(\gamma_9,a_9)I_9(\gamma_9,a_9)da_9\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_9(a_9) = (\rho_0)_9(0)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - (f_{90})_9 a_9}$$

$$I_9(\gamma_9, a_9) = \int_0^a (\rho_0)_9(\tau_9)e^{\gamma_9\tau_9 + (f_{90})_9\tau + \int_0^\tau \mu_9(s_9)ds_9}d\tau$$

$$= \frac{(\rho_0)_9(0)}{\gamma_9}(e^{\gamma_9 a_9} - 1)$$

Let

$$E_9(a_9) = e^{-(f_{90})_9 a_9 - \int_0^a \mu_9(\tau_9)d\tau_9}$$

Then $[Y_9]_9[\gamma_9]_9, a_9) = E_9(a_9)e^{-\gamma_9 a_9}$. by a simple calculation,

$$\int_0^\infty \beta_9(a_9)Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma}\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}\ da\right);$$

$$\int_0^\infty Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma_9}\left(\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}\ da_9\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_9(0)}{\gamma_9}[f_{90}^9]\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}da_9\ (\int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9)$$

$$+\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right).$$

$$\left(1+\frac{(\rho_0)_9(0)}{\gamma_9}(f_{90}^{'})_9\left(\int_0^{\infty}E_9(a_9)da_9-\int_0^{\infty}E_9(a_9)e^{-\gamma_9 a_9}da_9\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right)(1+\frac{(f_{90}^{'})_9}{\gamma_9}(\rho_0)_9(0)\int_0^{\infty}E_9(a_9)da_9)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (94), to

$$1-\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)da_9=1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\neq 0,$$

Since $\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)\int_0^{a}\alpha_9(\tau_9)Y^{-9}(\gamma_9,\tau_9)d\tau_9 da_9>0$. Therefore, the characteristic solution is

$$\gamma_9=-(f_{90}^{'})_9(\rho_0)_9\int_0^{\infty}E_9(a_9)da_9.$$

Since $f_9'(P_0)_9>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_9),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_9),0)$ which is related to the extinction of predators (debit summation types: that which acts as a resource)and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $( p^*(a), Q^* )$. It must satisfy the equations:

$$[p_a^*]_9 + (\mu_9(a_9) + [f_9]_9( P_9^*) + [f_9]_9(Q_9^*) )P_9^* = 0, \quad (26)$$

$$p_9^*(0) = h_9(Q_9^*) \int_0^\infty \beta_9(a_9)p_9^*(a_J)p_9^*(a_9)da_9 \quad (27)$$

$$g_9(P_9^*) - b_9 = 0, \quad (28)$$

Where $P_9^*=\int_0^\infty p_9^*(a_9)\, da_9$

Equation (28) yields $P^* = g^{-9}(b)$,

From (26),

$$p_9^*(a) = p_9^*(0)e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}$$

Substituting into (27), we need to solve

$$\left(1 - h_9(Q_9^*)\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}da_9\right)p_9^*(0) = 0$$

For $Q_9^*$

Fix $p_9^*(0)$ and $P_9^*$ Let

$$H_9(z_9) = 1 - h_9(z_9)\int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9}\, da_{,9}$$

And assume that

$$[k_3]_9 \int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(g_9^{-9}(b_9))a_9 - [k_9]_9 a_9}\, da_9 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_9(z_9)] > 0.$$

Clearly,

$$\frac{dH_9}{dz_9} = h_9(z)\left( [f_9^{'}]_9(z_9) \right)\int_0^\infty a_9\beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - f_9(P_9^*)a_9 - f_9(z_9)a_9} da_{\square}$$

$-h^{'}(z_9)\int_0^\infty \beta_9(a_9)\, e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9}\quad da_9$

y the assumptions on $\left[f_9^9\right]_9 and\ [h^9]^9 dH_9/dz > 0$. Hence, if

$$H_9(0) = 1 - \int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9} da_9 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_9(0) < 0$ follows if $g_9^-(b_9) < [f_9^{-9}]_9(r_9)$, where $r_9$ is the solution of (12).

Indeed, since

$$R_9(s_9) = \int_0^{\infty} \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 e^{-y\frac{a}{9}} da_9}$$

Is a decreasing function of $s_9$ when $P_9^* = [g^{-9}]_9(b_9) < [f_9^{-9}]_9(r_9)$, that is $[f_9]_9(P_9^*) < r$, $R\big(f_9(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_9(g^{-9}(b))) > 1, i.e.,$

$\int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9\left(g_9^{-9}(b)\right)a_9} > 1.$ (H4)

Then the system (4)-(6) has a unique positive equilibrium $(P_9^*(a_9), Q_9^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_9^*(a_9), Q_9^*)$ is more complicated. We again consider perturbations $m_9(t_9 a_9) = p_9(t_9, a_9) - p_9^*(a_9)$ and $W_9(t) = Q_9(t) - Q^*$.Then

$M_9(t) \equiv \int_0^{\infty} m_9(t, a_9) da_9 = P_9(t) - P_9^*,$

$$[m_9]_9 + [m_a]_9 = [p_9]_9 + [p_a]_9 - \left[p^*\right]_9(a_9) \approx -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9$$

$$-[f_{9*}^9]_9 p_9^*(a_9)M - [f_{9*}^9]_9 p_9^*(\mathrm{a})W_9,$$

$$m_9(t,0)p_9(t,0) - p_9^*(0) \approx h_* \int_0^{\infty} \beta_9(a_9) m_9(t, a_9)\, da + \frac{[h_*^9]_9}{[h_*]_9}\, p_9^*(0)W_9$$

$$\frac{dW_9}{dt} = \frac{dQ_9}{dt} \approx [g_*^9]_9 Q_9^* M_{,9}$$

Where $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $\left[f_{9*}^{'}\right]_9$ denoting $\left[f^{'}{}_9(P^*)\right]_9 \left[f_{9*}^{'}\right]_9$ denoting

$[f_9]_9$ $(Q_9^*,$ $h_9^*$ denoting $h_*^{'}$ denoting $h_9^{'}$ $(Q_9^*)$, and $\left[g_*^{'}\right]_9$

Denoting $\left[g^{'}\right]_9(P_9^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_9]_9 + [m_a]_9 = -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9 - \left[f^{'}{}_{9*}\right]_9 [p_9^*](a_9)M_9 - \left[f_{9*}^{'}\right]_9 p_9^*(a_9)W_9$$

$$m_9(t,0) = [h_*]_9 \int_0^\infty \beta_9(a_9)m_9(t,a_9)da_9 + \frac{[h_*^9]_9}{[h^*]_9}[p^*]_9(0)\, W_{,9} \qquad (30)$$

$$\frac{dW_9}{dt} = [g_*^9]_9[Q_9^*]M_9$$

Integrating $[m_9]_9(t,a_9)$ with respect to a from 0 to $\infty$, we have

$$\frac{dM_9(t)}{dt} = \int_0^\infty [mt]_9(t,a_9)da_9$$

$$= -\int_0^\infty [m_a]_9\,(t,a_9)da_9 - \int_0^\infty \mu(a_9)m(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + \left[f^{'}{}_{9*}\right]_9 \int_0^\infty p_9^*(a_9)da_9\right)M_9 - [f_{9*}^9]_9 \int_0^\infty p_9^*(a_9)da_9 W_9$$

$$= m_9(t,0) - \int_0^\infty \mu_9(a_9)\, m_9(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + \left[f_{9*}^{'}\right]_9 P_9^*\right)M_9 - \left[f_{9*}^{'}\right]_9 P_9^* W_{.9}$$

From (30), we arrive at the equations

$$\frac{dW_9(t)}{dt} = \int_0^{\infty} ([h_*]_9\beta_9(a_9) - \mu_9(a_9))m_9(t,a_9)da_9$$

$$-\left([f_{9*}] + [f_{9*}]_9 + \left[f_{9*}^{'}\right]_9 P_9^*\right) M_9 + \left(\frac{\left[h_*^{'}\right]_9}{[h_*]^9}\ P_9^*(0) - f_{9*}^{'} P_9^*\right) W_9,$$

(31)

$$\frac{dW_9(t)}{dt}\left[g_{*}^{'}\right]_9 Q_9^*\, M_{.9}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{"*}]_9\beta_9(a_9) - \mu_9(a_9))m_9(t,a_9)da_9 = ([h_*]_9\beta_9(\theta_9) - \mu_9(\theta_9))M_{,9}$$

Where $0< \ \theta_9 = \theta_9(t) < \infty$.

And if $a_9^9[a_9^9]2<4\ [a_9]2<4\ a_9e$,

$\text{Re}\,\lambda_9[\lambda_9]_9 = \text{Re}\lambda_9[\lambda_9]_9 = -\frac{9}{9}a_9[a2\ ]_9 < 0.$

Let $\Delta=\max_{9<u<\infty}\left|h_*[h_*]_9\beta_9(a_9) - \mu_9(a_9)\right|$. Then$|n_9(t)|$ ⍰⊿∀t ⍰0. Hence we obtain the following result.

**THEOREM 4**

Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_9^9 <4\ a_9e$ and $[a_9] > 1\Delta$,or $[a_9^9]$ ⍰4$[a_9]e$ and $(a_9 - \sqrt{(a_9^9)_9 - 4(a_9)_9e) > 1\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+9}\right)$ and $\beta_9(a_9) = ae^{-ma}$ Then $\Delta=[\mu_0] + 2$ Simple calculation leads to

$$\int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9) + d\tau_9 da} = \frac{1}{(m + \mu_0 + 1)^9} + \frac{1}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_9 + [\mu_0]$ 0.5. Assume$[f_9](x_9)= K_9 x_9$ Then the hypothesis (H9) is satisfied We take the conversion rate function $g_9(x_9)= [0x_9]/(2 + x_9)$; the functional responses $[[f_9](x) = [0]\, x_9/(1 + x_9)$and $h_9(x_9)=(1+\omega_9 x_9)/(1 + x_9)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_9 + [\mu_0] + [K][P_0])}{[m_9]_9 + (\mu_0)_9 + [K][P_0])_9^3\ [K]_9[P_0]_9^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_9$i e $[P_0] < b/(\delta_9 - b_9)$then the equilibrium$(P_0)$ $[P_0]_9(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_9$

$[f_{9*}] = K_9\ P^*\ [f_{9*}] = \frac{\sigma Q_9^*}{9+Q_9^*}$ $[h_*] = \frac{9+\omega_9 Q_9^*}{9+Q_9^*}$

$[f_{9*}] = K_9\ [f_{9*}^9] = \frac{\sigma_9}{(9+Q_9^*)^9}\ [g_*^9] = \frac{\sigma_9}{(9+P_9^*)^9}$ $[h_*^9] = \frac{\omega_9 - 9}{(9+Q_9^*)^9}$

In order to have (H3) and (H4) hold,$m, \mu_0, K, b, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_9}{(m+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9+b_9}+\sigma)^9}+\frac{1\omega_9}{([m_9]_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9+b_9}+\sigma_9)^3}<1$$

And

$$\frac{1}{(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9})^9}+\frac{1\omega_9}{(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9}+(m_9+[\mu_0]_9+\frac{K_9 b_9}{\delta_9-b_9}+\sigma_9)\,Q_9^*)^3$$

$$=m_9+[\mu_0]_9+3+\frac{K_9 b_9}{\delta_9-b_9}(1+Q^*)(1+\omega_9 Q_9^*$$

Now

$$\begin{bmatrix}1\\1\end{bmatrix}-4\begin{bmatrix}a\\9\end{bmatrix}e=\left(\frac{1K_9 b_9}{\delta_9-b_9}\right)^9+\frac{4K_9 b_9\sigma_9 Q_9^*}{(\delta_9-b_9)(9+Q_9^*)}+\frac{\sigma_9^9 Q_9^{*9}}{(1+Q_9^*)^9}$$

$$\frac{4\,b_9\sigma_9(\delta_9-b)Q^*}{\delta_9(1+Q_9^*)^9}-\frac{4\,(\delta_9-b_9)^9(1-\omega_9)Q_9^*}{\delta_9(1+\omega_9 Q_9^*)(1+Q_9^*)}P^*(0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_9 \, ae^{mu}_9 \quad [\mu_0]_9 - \frac{a_9}{a_9+9}$$

Hence if $[a^9_9 > 4(a_9)_9 e$ and

$$[a_9] = \frac{1K_9 b_9}{\delta_9 - b_9} + \frac{\sigma_9 Q^*_9}{1 + Q^*_9} > 1\Delta$$

Some of the examples to which the model has been applied are:

**New look at the galaxy-shaping power of black holes**

(358) "It creates (eb) these bubbles, and the bubbles keep the plasma hot.

(359) That's what prevents galaxies from (e) becoming even bigger than they are now."

(360) Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.

(361) In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.

(362) Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.

(363) "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.

(364) "It's like nature's thermostat, which keeps these galaxies from (e) growing.

(365) If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."

**A new look at the galaxy-shaping power of black holes**

(366) An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the

(367) McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.

(368) Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.

**Doomed satellite missions**

(369) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.

(370) At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.

(371) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.

(372) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.

(373) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.

(374) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people

(375) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."

(376) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.

(377) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(378) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.

(379) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.

(380) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.

(381) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.

| |
|---|
| (382) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well. |
| (383) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma. |
| (384) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6. |
| (385) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus. |
| (386) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (387) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (388) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (389) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (390) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (391) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (392) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (393) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (394) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (395) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (396) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (397) We know dark matter exists only from (e) its gravitational effects. |
| (398) It holds together galaxies and (e&eb) galaxy clusters. |
| (399) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (400) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |

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| (401) | They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (402) | X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (403) | They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (404) | This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (405) | An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |
| (406) | Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (407) | "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (408) | The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (409) | Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (410) | Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (411) | "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (412) | Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (413) | "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (414) | Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) |

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| consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories. |
| (415) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen. |
| (416) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer. |
| (417) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants. |
| (418) These were treated as (=) "aware" and "unaware" trials, respectively. |
| (419) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it. |
| (420) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another. |
| (421) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion. |
| (422) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness. |
| (423) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said. |
| (424) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness." |
| (425) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner. |
| (426) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain. |
| (427) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said. |
| (428) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

(429) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter.

(430) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined.

(431) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions.

(432) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.

(433) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.

(434) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.

(435) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.

(436) Noda and his colleagues instead studied the change in emissions over (e&eb) time.

(437) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.

(438) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).

(439) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.

(440) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.

(441) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson’s problem**

(442) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking “palm tree” organization.

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| (443) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society |
| (444) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible? |
| (445) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem. |
| (446) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other. |
| (447) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible. |
| (448) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges. |
| (449) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century. |
| (450) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching. |
| (451) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom. |
| (452) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth. |

**Spheres on trees**

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| (453) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration. |
| (454) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement. |
| (455) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum. |
| (456) By visualizing the problem in this way, the researchers noticed that these particular graphs don’t (e) have lots of branches extending from lots of (e) other branches. |
| (457) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure. |

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| (458) The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network. |
| (459) As a result, it provides (eb) an important clue for finding the global minimum. |
| (460) It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum. |
| (461) This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large. |

**Thompson's problem 2**

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| (462) By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society |
| (463) The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima. |
| **(464)** The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node. |

**Implications for Thomson's and Smale's problems**

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| (465) Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before. |
| (466) "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around." |
| (467) As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem. |
| (468) "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view." |
| (469) In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges. |
| (470) From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear. |
| (471) "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK. |
| (472) "For larger numbers, there will likely be (=) multiple funnels. |
| (473) Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum." |

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| (474) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem. |
| (475) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said. |
| (476) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org |

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| The nontrivial equilibrium is stable. |
| IF we define $[a_9]_9 = [f_{9*}]_9 + [f_{9*}]_9 + [f_{9*}]\,P^*_9$ |
| $$[a_9]_9 = \left[f^{9}_{9*}\right]P^*_9 - \frac{h^9_*}{h_*}P^*(0)$$ |
| $$e=[g^9_*]_9\,Q^*_9$$ |
| $$\eta_9(t) = \left[h_*\right]_9\beta_9\left(\theta_9(t)\right) - \mu_9\left(\theta(t)\right),$$ |
| The system(32)can be written as the following system: |
| $$[M(t)]\left[a_9\right]M_9 - [a_9]_9W_9 - [n(t)]M_9$$ |
| $$w_9(t) = eM_9$$ |
| Write |
| $$A=\begin{pmatrix} -[a_9]_9 - [a_9]_9 \\ e \qquad\quad 0 \end{pmatrix}$$ |

$$B=(t)=\begin{pmatrix} [\eta(t)]_9 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_9 \\ W_9 \end{pmatrix}$$

The solution of (32)is

$$X_9(t) = e^{A,t}[X]_9(0) + \int_0 e^{At}$$

**Note: There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and**

**sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by**

**term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation**

**alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation,**

**appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module**

**LEN NUNNEYS**

## Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

## Module numbered one

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems

and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \quad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1965); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be referred to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 6985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_6}{dt} = [N(t)]_6[(F_6)(R(t)]_6 - [M]_6[R(t)]_6$$

$$\frac{d[R(t)]_6}{dt} = [B]_6\,[(R)_6(t-T_6)] - [D]_6[(R(t)]_6 - [N(t)]_6\left[G\left(R(t)\right)_6\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_6}{dt} = a_6[N(t)]_6 - a_6[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_6 = a_3[N(t)]_3 - a_4[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_6 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

irst, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.6b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The

regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 6a is at the

**TIME DELAY**

FIG. 6. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.6$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**DELAY DUE TO PREDATOR MATURATION**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_6}{dt} = N_6(t - T_6)F_6(R_6(t - T_6) - N_6(t)M_6((R_6(t) \qquad (5)$$

$$\frac{dR_6(t)}{dt} = B_6\big(R_6(t) - D_6(R_6(t)\big) - N_6(t)G_6(R_6(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_6}{dR_6} + \frac{N_6 dG_6}{dR_6} - \frac{dB_6}{dR_6} > \frac{N_6 G_6}{2M_6}\left[\frac{dF_6}{dR_6} + \frac{dM_6}{dR_6}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_6/dR_6$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_6$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_6$) and the other derived from the non-delayed or "immediate" processes($I_6$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

ABSOLUTE STABILITY

$$L_6 = \frac{N_6 G_6}{2M_6}\frac{d(6nF_6)}{dR_6} (\geq 0) \qquad (6a)$$

$$I_6 = -\frac{N_6 G_6}{2M_6}\frac{d(6nM_6}{dR_6} (\geq 0) \qquad (6b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_6 > L_6 - I_6$, must be modified ahe $dM_6/dR_6$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_6 > L_6 - I_6, \quad \text{when} \quad P_6^2/4 > I \qquad (8a)$$

$$P_6 > L_6 - I_6 + \frac{(I_6 - P_6^2}{(L_6 - I)} \quad \text{when} \quad P_6^2/4 < I. \quad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

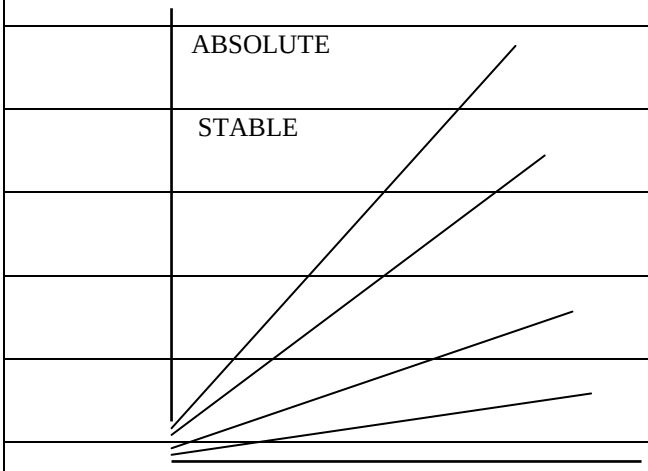



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

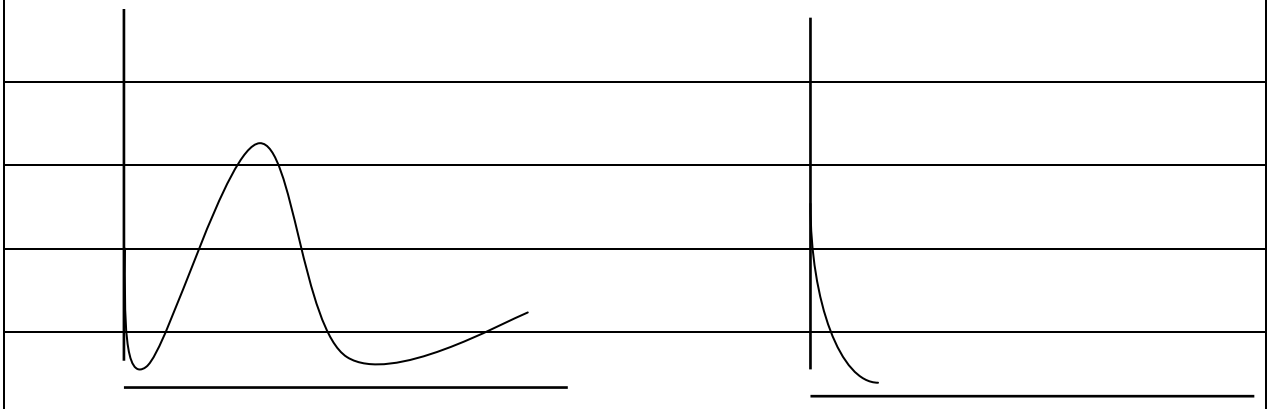
Can never completely offset the lagged effect,$L_6$; condition (8b) ensures that the total debit regulation, $P_6$, must always be positive even when $I_6$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_6$ and $L_6$. As $I_6$ increases relative to $L_6$ the zone of stability is enlarged; however,even when $I_6$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_6$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



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Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (6). Parameters in each case M=6; $dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6 = 0.05$ For (a) $NGdF/dR = 0.3, -N_6G_6dM_6/dR_6 = 0.6$; for (b)$NG_6dM_6/DR_6 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_6t}{dt} = N_6(t - T_6)F_6\big(R_6(t - T_6)\big) - N_6(t)M_6\big(R_6(t)\big)$$

(9)

$$\frac{dR_6t}{dt} = B_6\big(R_6(t - T_6)\big) - D_6\big(R_6(t)\big) - N_6(t)G_6\big(R_6(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) and the effective prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we

could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

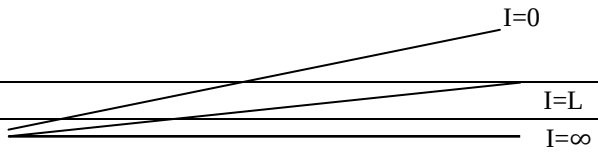
Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_6/dR_6$ is a measure of the control acting through the debit delay and setting $dB_6/dR_6 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_6/dR_6$ is a measure of the control acting through the credit delay although setting $dF_6/dR_6 = 0$ does not negate the effect of a predator delay; the value of $M_6$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_6/dR_6$and normalized by $M_6$ (see6a). The absolute stability of such a system is summarized in Fig. 6 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

TOTAL

ABSOLUTELY STABLE

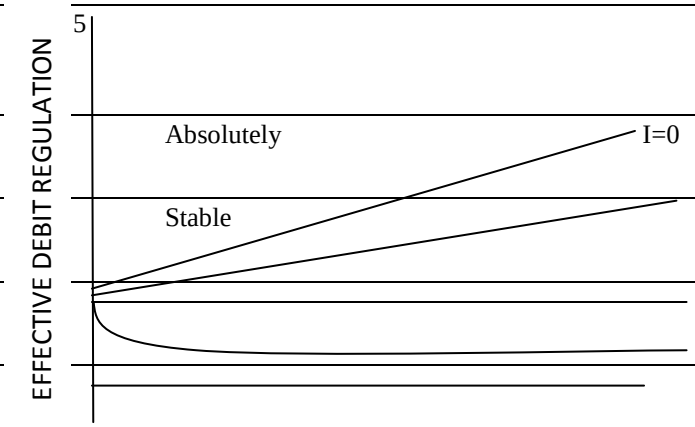



DELAYED DEBIT EFFECT

FIG6. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_6/dR_6 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.



-50 10

DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_6/dR_6 + N_6 dG_6/dR_6 + dB_6/dR_6)/M_6$). The immediate ($I_6$) and lagged ($L_6$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_6$(see 6b), is increased relative to the lagged part, $L_6$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_6 > I_6$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_6$is somewhat large than $I_6$ then the condition constraining prey regulation of recruitment ($dB_6/dR_6 < 0$) is

$$\frac{dD_6}{dR_6} + N_6\frac{dG_6}{dR_6} + \frac{dB_6}{dR_6} > \frac{(L_6 - I_6)M_6}{2} \qquad (10)$$

If $L_6 < I_6$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_6$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_6$

The - $D_6$ positions are defined by letting $\lambda_6 = iZ_6$ and separating (A1.2) into real and imaginary parts:

$$Z_6^2 + V_6 Z_6 \int_0^{\infty} \sin(Z_6 T_6) f(T_6) dT_6 - (X_6 + Y_6) = 0$$

$$Z_6 \left( U_6 - V_6 \int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6 \right) = 0$$

Noting that $X_6 + Y_6 = 0$ is we assume X+Y=0.It follows that $Z_6 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_6 = V_6 \int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6$$

If $U_6 > |V_6|$then given the constraints of $f(T_6)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A6.5)remains

greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

## ABSOLUTE STABILITY WITH A DELAY DUE TO CREDIT CONSOLIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_6^2 + (U_6 - V_6 + M_6)\lambda_6 - M_6\lambda_6 e^{-T_6\lambda_6} + [(U_6 - V_6)M_6 + Y_6] - [((U_6 - V_6 M_6 -]e^{-T_6\lambda_6} = 0$$

(For definitions see Appendix6) We follow the technique outlined in Appendix 6 and let $\lambda_6 = iZ_6$ and separate (A2.6) into real and imaginary parts

$$Z_6^2 + M_6 Z_6 \sin(Z_6 T_6) - (A_6 M_6 + Y_6) + (A_6 M_6 + X_6)\cos(Z_6 T_6) = 0$$

$$(A_6 M_6) Z_6 - M_6 Z_6 \cos(Z_6 T_6) + (A_6 M_6 - X_6)\sin(Z_6 T_6) = 0$$

Where $A_6 = U_6 - V_6$Using (A2.3), substitute for $Z_6^2$ and $Z_6$in (A2.2):

$$W_6^2(X_6^2 - 3M_6 A_6 X_6 - M_6^2(X_6 - Y_6) + 2M_6^3 A_6 + 2A_6^2 M_6^2]$$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_6 = 1 - \cos(Z_6 T_6)$.

In the absence of a delay the system is stable for $A_6 > 0,$ and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_6 > (X_6 - Y_6)/2\ M_6$ Thus a necessary condition for absolute stability is

$$A_6 > [\ X_6 - Y_6/2\ M_6$$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_6$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_6$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_6, A_6) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \leq W_6 \leq 2$ *is* of interest First consider the extreme values

$$Q_6(0, A_6) = A_6^2\ (X_6 - Y_6)$$

$$Q_6(2, A_6) = (A_6 + 2\ M_6)^2 (2A_6 M_6 - (X_6 - Y_6))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_6}{\partial A_6} = W_6^2\left(2M_6^3 + 4A_6 M_6^2 - 3M_6 X_6\right) + W_6\left(3M_6 A_6^2 - 2A_6 X_6 + 2M_6\left(2X_6 + Y_6\right)\right)$$

$$+2A_6\left(X_6+Y_6\right)$$

$=2M_6^2W_6^2(M_6+A_6)+M_6A_6^2W_6+\left(2A_6M_6-X_6+Y_6\right)(A_6W_6+M_6W_6^2$

$+(2-W_6)((2M_6X_6+M_6Y_6)W_6+A_6(X_6+Y_6)),$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_6$quadratic in $W_6$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_6$has no minimum in the range 0≤ $W_6$ ≤2.Defining

$Q_6\left(W_6,A_6\right)=a_6W_6^2+c_6 where\ a_6\ b_6\ and\ c_6$ Are functions of $A_6$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_6≥b_6≥0$$

And

$$b_6^2≤4a_6c_6$$

$$Case(i):Y_6=0$$

$Let\ Y_6 = 0$ And try the substitution A= $X_6 = /2M_6\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_6 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_6 X_6 /$ $2M_6$) is positive for *W<2. Since* $Q_6$ increases with $A_6$(see A2.9), increasing $A_6$will generate no further solutions of $Q_6$=0. Hence the condition A>X/2M defines absolute stability when $Y_6 = 0$.

$Case\ (ii)$: $Y_6 = y_6\ X_6$ ,

$Let\ Y_6 = y_6\ X_6$And try the substitution $A_6 = X_6/(2 + A_6\ y_6)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_6/(2 + n_6 y_6)$ we have

$$W_6^2\ y_6\left[X_6 n_6 \frac{(1 + n_6 y_6)}{(2 + n_6 y_6)} + M_6^2\left(2 - n_6 + n_6 y_6\right)\right]$$

$$-W_6\ X_6\left[\frac{X_6(1 + n_6 y_6)}{M_6^2(2 + n_6 y_6)^2} + 2\ y_6\left(n_6 - 1\right) + \frac{X_6^2(6 + y_6)}{M_6^2(2 + n_6 y_6)} = 0.\ (A2.12)\right.$$

Since $n_6$ and $A_6$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_6$

Defining $E_6 = (2 + n_6 y_6), the\ condition (A2.66?) can\ be\ written\ as$

$[\frac{X_6(E_6 - 6)}{E_6^2}]^2[\frac{X_6^2}{M_6^2} - 4\ y_6\ E_6^2\frac{X_6}{M_6^2} + 4y_6(n_6 - 2)E_6^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in ($X_6 / W_6^2$) can be expressed as

$2E_6[y_6E_6+(y_6E_6(6+y_6)(E_6-n_6))^{6/2}]>\frac{X_6}{M_6^2}>2E_6[y_6E_6-y_6E_6(6+y_6)(E_6-n_6))^{6/2}]$

**ABSOLUTE STABILITY**

Hence if $X_6/M_6^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0$W_6$2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_6-6)}{E_6^2}\left(\frac{X_6}{M_6^2}\right)^2-2y_6\frac{(E_6+n_6^2y_6)}{E_6}\left(\frac{X_6}{M_6^2}\right)-4y_6(E_6-n_6)$$0

And has the solution

$\frac{X_6}{M_6^2}>\frac{E_6}{(E_6-6)}[y_6(E_6+n_6^2y_6)+[y_6^2(E_6+n_6^2y_6)^2+4y_6(E_6-6)(E_6-n_6)]^{6/2}]$

Given $E_6>n_6$this lower limit to $X_6/M_6^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_6/M_6^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_6$in the appropriate range and this limit is further lowered by (A2.64)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_6<n_6(ie.,2A_6M_6^2<X_6-Y_6$)then (A2.6)and

(A2.8)show that $Q_6$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_6}{M_6^2} > 2E_6[y_6E_6 - (y_6E_6(6 + y_6)(E_6 - n_6)]^{6/2}]$$

This equations throws light on the individual general ledger and Nature's general ledger as we call it : see various papers on the issue. powers, which are the energies of Matrka we have already discussed, are not the only obstacles the yogi must overcome. He must, for example, also resist the temptation to rest content with the miraculous yogic powers (siddhi) he acquires in the course of his spiritual develop- ment. Again to do this he must practice Yoga. Similarly, in order to pervade the Fourth State gradually through the other states in the manner proper to practice at the Individual level, the method is the same. He must practise the higher yoga of the Tantras which, turning his mind inwards and freeing it from discursive representations, allows him to penetrate into the Supreme Principle. 312 Once the yogi has attained this contemplative state, his main problem is to make it permanent. In the introverted state the gross external movement of the breath is suspended and with it the activity of the intellect, mind, individualised consciousness, powers of the senses and the ego. 313 When the yogi rises out of this state, he is liable to fall again into the lower order of creation generated by Maya if he does not maintain his awareness of the higher reality he has experienced and allows his awakened, illumined insight to be obscured by the dream-like vision of thought-constructs. 314 Naturally, the yogi must rise out of the introverted condition of suspension. It is inherent in the very nature of reality that it should move out of itself. 315 Pure, universal consciousness initially transforms itself into the vital breath 316 charged with the impression iyasana) of the power of awareness attained through introversion. By attending to the pulse (spanda) of the breath as it moves out of the absolute, the

yogi can develop an intuitive sense of the inherent unity of all he will perceive in the mental and physical spheres created by the outpouring of consciousness. In this way he realises that his own nature is everywhere present in all he perceives and that all things thus reside within him. Blessed with this insight his consciousness remains free and unlimited even at the individual level where the breath, mind, senses and body are active. If the yogi fails to do this, he finds himself once again beset by the strictures of his embodied existence and must, as before, try to pervade all his other states of consciousness with the aesthetic delight (rasa) and wonder (camatkara) of the Fourth State he experienced in contemplation. Again this means that he must strengthen his pure, empowered awareness that his universal nature manifests as all things. 317 In this way he discovers Siva's presence in every sphere of individualised consciousness ranging from the breath to outer objectivity. The yogi's mind then becomes tranquil and undistracted because wherever it may wander, the yogi perceives only Siva, his authentic nature. 318 Consciousness is thus freed of all external referents and the yogi's subjectivity is purified of all identification with the body or anything else that belongs to the objective sphere. The yogi then becomes detached from the opposites of pleasure and pain and is transcendentally free (kevalin). 319 yogis are again, however, liable to fall if he allows himself to get entangled in the play of opposites. This fall is more serious than the others because, although he is caught by the confining restrictions of individual- ised consciousness as before, he is now also affected by karma. Fleeing from pain in the pursuit of pleasure he is bound to act (karma) to minimise one and maximise the other and so is thrown down to the lowest level of embodied subjectivity (sakala). In order to regain his lost state, he must ascend gradually, by Siva's grace, from one order of subjectivity to the next and so free himself progressively of the limitations of the lower levels to gain the greater freedom and expansion of the higher. As he progresses, the objective sphere also evolves from the grossest perceptions of physical objects outside the lowest order to subjectivity, through to the subtler inner, mental

perceptions to finally reach the order of subjectivity that contains objectivity within itself and is free to externalise it at will. 320 The degree to which this process develops depends, as before, on the yogi's awareness of the Fourth State. In consonance with the general principle that the remedy should suit the defect, the yogi is instructed to seek this higher state of consciousness in the wonder (camatkara) or delight (ananda) he feels in moments of intense physical pleasure. At first he experiences this subtle consciousness for an instant in the subjective sphere. If he manages to catch hold of it, it becomes more intense as the cognitive and objective spheres are also gradually pervaded and vitalised by it. Occasions for this practice are, for example, the sense of satisfaction one feels after a good meal or the aesthetic delight one experiences when listening to good music or the pleasure of sexual union with the Tantric consort or even solitary sexual excitation. In these moments of delight the yogi can penetrate momentarily into his own authentic Siva-nature (sambhavavesa) through the empowered contact (saktasparsa) m he makes with it in the freedom of the pure subjectivity of the Fourth State. 322 If the yogi develops his awareness of this higher level of consciousness and maintains it, he eventually experiences it constantly. 323 Clearly, what prevents the yogi from attending to his state of consciousness rather than the circumstances which induce it is the craving for pleasure (abhila§a) born of ignorance — the source of every impurity which clouds consciousness. Craving directs the yogi's attention towards outer, worldly things and so he is caught in the net of thought- constructs. 324 To free himself of his worldly desires and reverse this binding extroversion, the yogi must eradicate its cause. To be freed of all the ups and downs of the path and no longer be tormented by the possibility of a fall, the yogi must see reality perfectly and completely. This insight is itself liberation and the moment it dawns the yogi is instantly freed. This sudden realisation is the goal of Tantric Yoga. E Tantra declares: "He who perceives reality directly, even for the brief moment it takes to blink, is liberated that very instant and never reborn again." 325 Although the yogi's body and mind

continue to function as before, they are like mere outer coverings 326 which contain, but do not obscure, the mighty, universal consciousness which operates through them. The yogi's body is (=) the universe, the senses (=) the energies that (eb, eb+)vitalise it, his mind (=) Mantra, the rhythm of his breath (eb) the pulse of time and his inner nature pure, dynamic consciousness. Raised above all practice, and hence all possibility of falling to lower levels, the yogi realises that he has always been free 327 and that his journey through (e&eb) the **dark land of Maya was nothing but a dream, a construct of his own imagination** The channels of the senses by which the yogi enjoys the pleasure of the objects of sense are those through which he fills the three worlds with the consciousness emitted from his own Heart. 78 **The Doctrine of Vibration T**he senses are outer, subsidiary cycles of conscious energy (anucakra), revolving around the core or primary cycle (mukhyacakra) of pure consciousness. They scintillate with the diversity of sensations and resonate in the field of their awareness like the feathers of a peacock unfurled and folded back in the ecstasy of its dance. 79 Ultimately it is Siva Himself Who, for His own pleasure, 80 extends and withdraws them, 81 not the individual soul. The outward movement of the senses is the creative flow of consciousness from the subject who, like an ocean, is the source of the waves of their activity and the final resting place of the flow of the rivers of sensations and thought-forms that unfold through the arising of the power of each of the senses. 82 The Stanzas on Vibration declare: Indeed the individual soul does not activate the impulse of the will [which directs the body's activity] by himself alone but through his contact with [his] own [inner] strength {bald) made in such a way that he identifies with it [thus acquiring its power]. 83 The close contact between the senses and universal consciousness vitalises them and sets them in motion. They share in the properties of consciousness, just as a heated iron ball can burn, cook or give out light. 84 Spanda doctrines does not agree with the Sarnkhya view that the individual soul (purusa) alone stimulates activity in the objective sphere (prakrti) by the contact he makes with it. The soul can do

this only because universal consciousness animates him. Moreover, consciousness, which has the power to impel this activity, cannot itself be inactive as the Sarnkhya maintains. The influence it has on lifeless matter implies an active exertion. 85 This activity is the subtle, inner outpouring of the Self (svatmocchalatta) within the Self. Outside the confines of time and space, in the centre of the circle of the senses, this inner activity initiates the expansion (prasara) of the flow of conscious energy. As the pure exertion (udyoga, udyama) of Being (bhava), it is the source of the initial unfolding (unmesa) of Becoming. 86 Ksemaraja explains: [By] 'exertion' (udyama) is [meant] the emergence of the supreme [level of] intuition (pratibha) which is the sudden outpouring of the [immanent and transcendental] unfolding of the reflective awareness of pure consciousness. 87 The direct, intuition (pratibha) of reality, free of thought-constructs, is the supreme power (parasakti) of consciousness. The yogi can gain Divine **Body and Sacred Circle of the Senses 153** this pure intuitive awareness by grasping the initial expansion (unmesa) of consciousness which engenders the emergence of thought and the activity of the senses. It can be discovered in the subtle Centre, between one thought and the next. The Stanzas on Vibration teach: The expansion of consciousness that takes place when one is engaged in a single thought should be known to be the source from whence another arises. One should experience that for oneself! 88 This does not mean that the unfolding expansion of consciousness {unmesa) is an intermittent, transitory occurrence. 89 As the source of all thought, including the notions of past, present and future, 90 it stands outside the confines of time. Consciousness is in a state of perpetual expansion. It is an ever developing process, a tiny part of which we observe as the flow of events around us and the thoughts and feelings within us. The totality of this process can never be grasped objectively. Objectivity is a state of fragmentation (bheda) in which discrete elements are discernible and mutually distinguished as individual, specifically definable entities. It is the realm within which thought and language operate. To grasp reality in its completeness, we must go

beyond the partial representations of thought and speech. To experience the primordial source and basis of all things, we must pierce through the outer periphery of thought and plunge into the Centre, to discover the instant in which thought, and with it, the sense of diversity, initially emerges. Ksemaraja cites a practice from the Vijnanabhairava through which the yogi can achieve this state of awareness: Checking [the movement of] attention (cit) once it has quit its object, it must not [be allowed] to move to any other; contemplation (bhavana) then blossoms forth by [experiencing] the state in the centre between them (tanmadhyabhava). 91 Left unattended and undisciplined, awareness spontaneously shifts from one object of thought to (e&eb) the next. The transition from one to the other entails a movement through (e&eb) a state of pure indeterminate (nirvikalpa) awareness divested of all thought-forms and perceptions. The yogi who succeeds in checking the movement of his attention, experiences within himself the subsidence of thought into the emptiness of the Centre which unfolds full of the powerful pulsation of consciousness within which all objectivity merges and becomes one 92 in the trans- cendental outpouring (lokottarollasa) of its aesthetic rapture. 93 turning in on himself, the yogi must exert his power of awareness to (e) discern the vitality of (e) his own blissful Spanda nature. Just as a distant object becomes (=) more clearly visible by (e) exerting an effort to (e) perceive it, in the same way, through (e&eb) practice, the yogi's own Spanda nature becomes (=) progressively more evident to him, 94 The effort the yogi exerts to free himself of ignorance is essentially one with the force of awareness that impels the senses and mind. This inner exertion is the basis of every means to liberation. Ultimately bestowing the bliss of self- realisation, it carries the yogi along the path to salvation with all its twists and turns like wind which blows a cloth in various directions. 95 As the basis and ultimate goal of all yogic practice, this exertion is the pure wonder (camatkdra) inspired by the realisation of the essential identity of consciousness (cit) and being (sat). 96 As such it is the divine radiance (sphuratta) of the Heart of one's own consciousness 97 where the powers of

will, knowledge and action merge in the harmony (samarasya) of bliss. 98 Thus the highest form of worship (varivasya) is the contemplation of the inner strength (bald) or universal exertion of one's own true nature. 99 "One should examine [that reality]," says Ksemaraja, "with the exertion which is Bhairava's nature and is true service (sevana) to the integral inner nature [of all things] and the unfolding of the activity of one's own vitality (ojas) which withdraws all duality [back into consciousness]." 100 Mindful of the eternal joy (ahlada) of his inner strength, the yogi must give up all desire for the fleeting moments of petty pleasure (ksanikasukhalava) ]0] he may glean from outer objects. He must even give up the desire to achieve enlightenment and a reality (tattva) which can never be realised through a personal act of will. 102 In short, he must die to himself in an act of profound faith (sraddha) and adoration (adara) to find rest in the supreme reality free from the binding dichotomy of means and objectives (upeyopayabhava) — the reality of his perfected effort. The yogi, now full of devotion for this reality unfolding within himself and the universe (antarmukhatattva), merges into it and, by achieving the pure consciousness and bliss (cidananda) of the power (bala) of his own nature, achieves oneness. 103 The yogi is firmly established in his true nature by recognising the creative unfolding (unmesa) of consciousness through which he is able to know all things, whether in the present, past or future. 104 Master of the vitality which gives life to his mind and body, he can strengthenit directly without need of food or drink, 105 or quit it at will to vivify Divine Body and Sacred Circle of the Senses 155 that of another. 106 Experiencing the eternal satisfaction of his own uncreated nature, 107 the yogi is freed of disease and the infirmities of old age. 108 Inspired by this vitality, he has the inner strength to overcome the severest hardships and obstacles in the way of his ultimate goal. 109 The Stanzas on Vibration explain: Lassitude ravages the body and this arises from ignorance but if this is eliminated by an expansion of consciousness (unmesa), how can [ignorance], deprived of its cause, continue to exist? 110 Doubt (sanka) is the source of every spiritual ailment. 111 A

man in doubt about his true nature, and the path to follow by which he may come to realise it, is constantly overcome by difficulties. 112 "Due to ignorance" teaches the Sarvacaratantra, "the deluded man is in doubt and thus [suffers the cycle] of birth and death." 113 Caught between conflicting alternatives, the mind shifts from one to the other becoming more entangled in its own thoughts as it does so. Doubt thus contracts consciousness 114 sullying it with the turbulence of confused thinking. 115 At once both the root of "the ancient tree of transmigration" and the first sprout of its seed, 116 doubts deprives man of the innate bliss (sahajananda) of his own nature. 117 Like a thief it steals away the wealth of true knowledge and reduces man to an imagined state of poverty in which he feels hemmed in on all sides by constraints and limitations. 118 Ignorance and its sister, doubt, are the essence of all impurities (mala) which sully consciousness. To the degree in which the yogi is cured of this ailment, namely, the lassitude (glani) of doubt, his true nature becomes manifest, just as raw gold when heated is freed from dross. When the yogi recognises the all-powerful expansion (unmesa) of his consciousness, he is projected beyond the realm of relative distinctions between virtue and vice. Thus overcoming all doubt as to what he should or should not do, he penetrates into the pulsing Heart of Bhairava's consciousness. 119 As the Stanzas on Vibration say: An individual who [though] desirous of doing various things [but] is incapable of doing them due to his innate impurity, [experiences] the Supreme state when the disruption (ksobha) [of his ego] ceases. Then [the soul realises] that his [true] uncreated nature is [universal] agency and perceiving subjectivity and so he knows and does whatever [he] desires. 120 the individual soul is(=) a combination of (e) conscious (cit) and 156 The Doctrine of Vibration unconscious (acii) elements 121 brought together when he allows the power of his authentic nature to be obscured by (e) his own impurity. This impurity is (=) three-fold. The first is technically called (=) the 'Impurity of Individuality' (anavamald). Due to (e) this impurity, the individual soul fails to (e) recognise his all-embracing fullness (purnata) and is disturbed by (e) his craving

for (e) experience (bhoga) based on (e) the unreasoned assumption that he is (=) incomplete (apurnd) and in need of (e) something outside himself. 122 Thus it manifests as (=) desire or attachment (rdga). ]23 At the same time it is (=) primordial ignorance, 'the cause of the sprout of (e) transmigration'. 124 As such, it works in two ways. Firstly, it deprives (e) individual consciousness of freedom. The yogi affected by (e&eb) this impurity is incapable of (e) entering or emerging out of the states of contemplation as he wishes. Secondly, this impurity renders (eb) the individual unconscious of his freedom, which for(e) the yogi means that he is (=) liable to lose consciousness of (e) his true nature when he is (=) not in a state of (e) contemplative absorption. 125 The second impurity is (=) called the 'Impurity of Maya' (maylyamald). This comes into operation when consciousness has been contracted by (e) the Impurity of Individuality. It limits (e) the power of knowledge operating in (e&eb) the individual soul through (e&eb) the inner and outer organs of (e) perception. 126 Due to this impurity he perceives (e) diversity everywhere. 127 Finally, the third impurity, that of Karma, comes into play when (e) the individual, deprived of (e) the freedom and knowledge of (e) unconditioned consciousness, acts (e&eb) in his limited way prompted by (e) desires and fears for (e) his personal gain. Disturbed by (e) these three impurities the soul cannot find rest in (eb) himself.

And

$$E_6 > n_6$$

Replacing $E_6$ and $n_6$ in these conditions leads to the form

If $A_6^2 > 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6$.

If $A_6^2 < 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6 .+ (Y_6 - A_6^2/2)^2/2M_6(X_6 - Y_6)$.

Note that by normalizing $A_6$ by $M_6$ and $X_6$ and $Y_6$ by $W_6^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_6^2 + (M_6 + U_6)\lambda_6 - (M_6 + V_6)\lambda_6 e^{2T_6\lambda_6} + (M_6 U_6 + Y_6)$$

$$-[M_6(U_6 + V_6) - X_6]\, e^{2T_6\lambda_6} + M_6 V_6 e^{2T_6\lambda_6} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_6 = \mathrm{i} Z_6$ *and* separate

(A3.1)into real and imaginary parts

$$Z_6^2 + (M_6 + V_6) Z_6 \sin(Z_6 T_6) - (U_6 M_6 + Y_6) + (M_6(U_6 + V_6) - X_6)$$

$$\mathrm{X}\cos(Z_6 T_6) - 2M_6 V_6 \cos^2(Z_6 T_6) + M_6 V_6 = 0$$

$$(M_6 + U_6) Z_6 - (M_6 + V_6) Z_6 \cos(Z_6 T_6) + (M_6(U_6 + V_6) - X)$$

$$\mathrm{X}\sin((Z_6 T_6) - M_6 V_6 \sin(Z_6 T_6) = 0$$

Using (A.3.3) to eliminate $Z_6$ and $Z_6^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_6 = 1 - \cos Z_6 T_6$ and $A_6 = U_6 - V_6$ we obtain

$$4 M_6^2 V_6^2 W_6^4 + 2 M_6 V_6 W_6^3 [(M_6 - V_6)^2 + A_6(M_6 + V_6) + 2(A_6 M_6 - X_6)]$$

$$+ W_6^2 [M_6 - V_6)^2 (2 A_6 M_6 + Y_6 - X_6) + (2A_6 M_6 - X_6)^2 - A_6 (M_6 - V_6)(2A_6 M_6 - X_6)$$

$$+ M_6 V_6 (A_6^2 + 4(X_6 + Y_6))]$$

$$+ W_6 [A_6^2 + 2X_6)(2 A_6 M_6 - X_6) + 2A_6 Y_6 (M_6 + V_6) - M_6 A_6^3] + A_6^2 (X_6 M_6 - Y_6) = 0$$

Being a quartic in $W_6$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_6 > 0$$

$$and$$

$$C_6 > (X_6 - Y_6)/2 M_6$$

$recalling\ that\ A_6 = U_6$ and defining $C_6 = U_6 + V_6$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_6 (0, U_6 V_6 = A_6^2 (X_6 + Y_6)$$

$$Q_6 (2, U_6 V_6) = (C_6 + 2M_6^2)^2 (2C_6 M_6 + Y_6 - X_6)$$

So that given (A3.6) the extreme $W_6$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_6$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_6$ the existence of a D-position at this point is

determined by letting cos $(Z_6 T_6) = -1\ in\ (A3.3)$

ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 e^{-\tau_6\lambda_6} + X_6 + Y_6 = 0 \qquad (A1.1)$$

Where $V_6 = dB_6/dR_6$ $U_6 = dD_6/dR_6 + N_6 dG_6/dR_6$ $X_6 = N_6G_6dF_6/dR_6$ and $Y_6 = -N_6G_6dM_6/dR_6$ (evaluated at the equilibrium). Note that $U_6X_6$ and $Y_6$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_6$ as $(T_6)$ . (It follows that $f(T_6)$ 0 for all $T$, and $\int_0^\infty f(T_6)dt = 1.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 \int_0^\infty e^{-\tau_6\lambda_6} f(T_6) + X_6 + Y_6 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MODULE NUMBERED TWO**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed

the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).. Haimovici's and Nunny's models have been applied to some of the Deleuzean and Heideggerian problems too : Scoring on that count are :

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time

delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t - T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if

only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

TIME DELAY

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

## MATURATION AND CONSUMMATION OF CONFEDERATION OF DEBIT AND CREDIT TYPES:

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t-T_7)F_7(R_7(t-T_7) - N_7(t)M_7((R_7(t) \quad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7} + \frac{N_7 dG_7}{dR_7} - \frac{dB_7}{dR_7} > \frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7} + \frac{dM_7}{dR_7}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes $(L_7)$ and the other derived from the non-delayed or "immediate" processes$(I_7)$:as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_7 = \frac{N_7 G_7}{2M_7}\frac{d(7nF_7)}{dR_7} (\geq 0) \quad (7a)$$

$$I_7 = -\frac{N_7 G_7}{2M_7}\frac{d(7nM_7}{dR_7} \ (\geq 0) \qquad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7 > L_7 - I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_7 > L_7 - I_7, \quad \text{when} \quad P_7^2/4 > I \qquad (8a)$$

$$P_7 > L_7 - I_7 + \frac{(I_7 - P_7^2}{(L_7 - I)} \quad \text{when} \quad P_7^2/4 < I. \qquad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

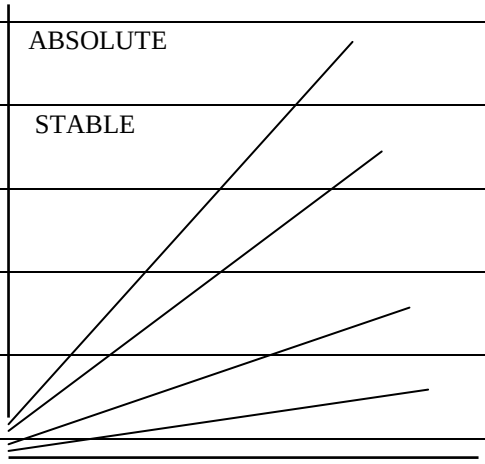



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit

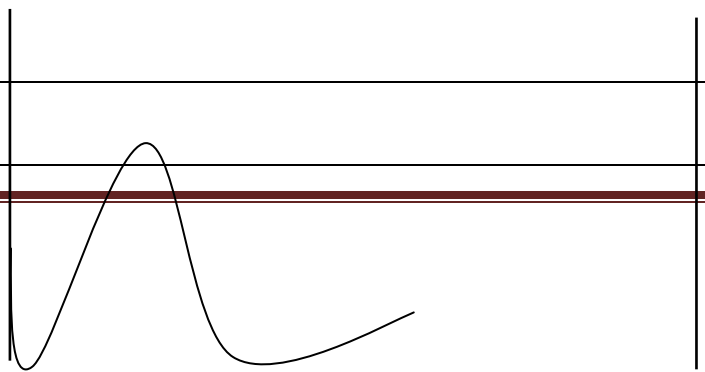
(like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

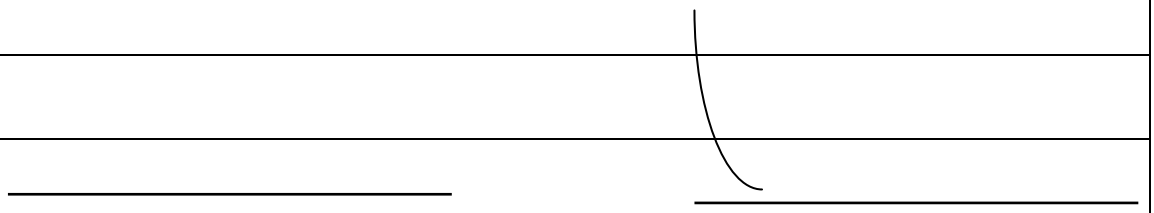
It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the





TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t - T_7)F_7\big(R_7(t - T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t - T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation $(dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7)$ and the effective prey regulation $(dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7)$ must be positive.

The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the

total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

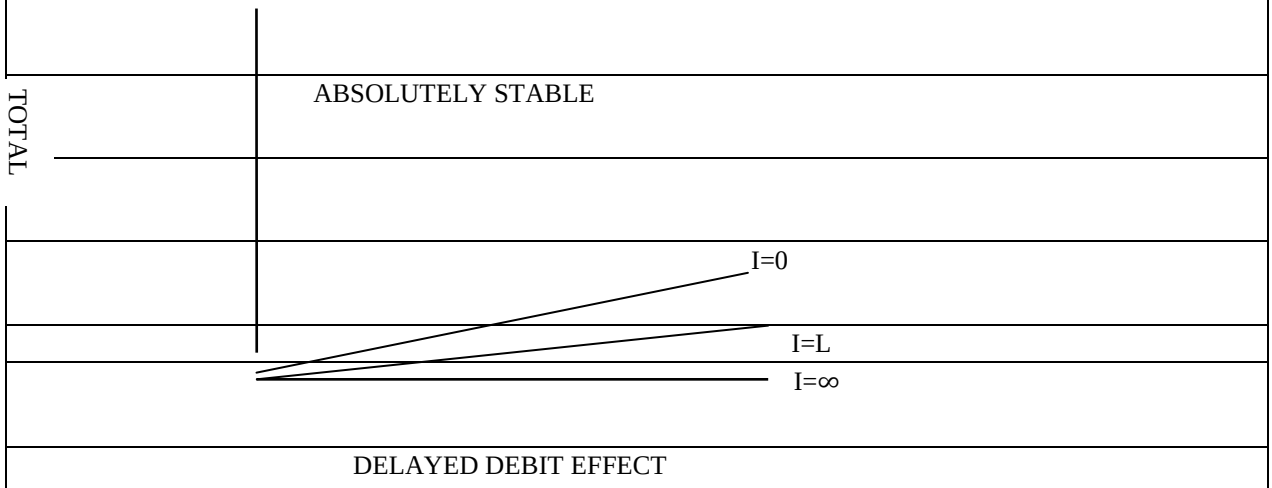
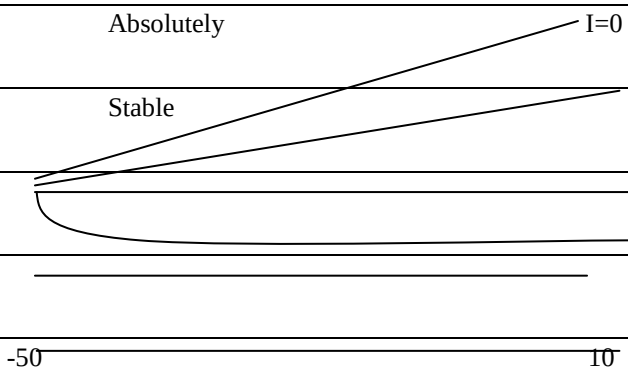



FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay.

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^\infty \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7\left(U_7 - V_7\int_0^\infty cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7\right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7\int_0^\infty cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$then given the constraints of $f(T_7)$, $(A1.5)$*can never be satisfied Co*nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7\lambda_7 e^{-T_7\lambda_7} + [(U_7 - V_7)M_7 + Y_7] - [((U_7 -$$

V 7M 7−]**e**−7 7λ 7=0

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$\boldsymbol{Z}_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7)\cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7)\sin(Z_7 T_7) = 0$$

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $\boldsymbol{Z}_7^2$ and $Z_7$in (A2.2):

$W_7^2(X_7^2$-3 $M_7\ A_7\ X_7$- $M_7^2(X_7$- $Y_7) + 2\ M_7^3\ A_7$+2 $A_7^2\ M_7^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_7$=1-cos $(Z_7 T_7)$.

In the absence of a delay the system is stable for $A_7 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7 > (X_7 - Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$$A_7 > [\ X_7 - Y_7/2\ M_7$$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_7, A_7) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_7 \le 2$ *is* of interest First consider the extreme values

$$Q_7(0, A_7) = A_7^2\ (X_7 - Y_7)$$

$$Q_7(2, A_7) = (A_7 + 2M_7)^2 (2A_7 M_7 - (X_7 - Y_7))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7 M_7^2 - 3M_7 X_7\right) + W_7\left(3M_7 A_7^2 - 2A_7 X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7+Y_7\right)$$

$=2M_7^2W_7^2(M_7+\mathrm{A}_7)+M_7A_7^2W_7+\left(2A_7M_7-X_7+Y_7\right)(A_7W_7+M_7W_7^2$

$+(2-W_7)((2M_7X_7+M_7Y_7)W_7+A_7(X_7+Y_7)),$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $W_7$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$has no minimum in the range $0 \le W_7 \le 2$.Defining

$Q_7\left(W_7,A_7\right)=a_7W_7^2+c_7 where\ a_7\ b_7\ and\ c_7$ Are functions of $A_7$we can define the condition for a minimum in the range 0-2

condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7 \ge b_7 \ge 0$$

And

$$b_7^2 \le 4a_7c_7$$

$$Case(i): Y_7=0$$

$Let\ Y_7=0$ And try the substitution A= $X_7=/2M_7\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}\ (W_7 X_7 / 2M_7)$ is positive for *W<2. Since* $\mathrm{Q}_7$ increases with $A_7$(see A2.9), increasing $A_7$will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

$Case\ (ii): Y_7 = y_7\ X_7$ ,

*Let* $Y_7 = y_7\ X_7$And try the substitution $A_7 = X_7/(2 + A_7\ y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + n_7 y_7)$ we have

$$W_7^2\ y_7\left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2\left(2 - n_7 + n_7 y_7\right)\right]$$

$$-W_7\ X_7\left[\frac{X_7(1 + n_7 y_7)}{M_7^2(2 + n_7 y_7)^2} + 2\ y_7\left(n_7 - 1\right) + \frac{X_7^2(7 + y_7)}{M_7^2(2 + n_7 y_7)}\right] = 0.\ (A2.12)$$

Since $n_7$ and $A_7$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7), the\ condition (A2.77?) can\ be\ written\ as$

$\left[\frac{X_7(E_7 - 7)}{E_7^2}\right]^2\left[\frac{X_7^2}{M_7^2} - 4\ y_7\ E_7^2\frac{X_7}{M_7^2} + 4y_7(n_7 - 2)E_7^3\right]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$2\ E_7[y_7 E_7 + (y_7\ E_7(7 + y_7)(E_7 - n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - y_7 E_7(7 + y_7)(E_7 - n_7))^{7/2}]$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2$) is large or small then (A2.6) will have a real solution although if may not be in the range (0 W 7 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7 - 7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2 - 2y_7\frac{(E_7 + n_7^2 y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right) - 4y_7(E_7 - n_7)$$ 0

And has the solution

$\frac{X_7}{M_7^2} > \frac{E_7}{(E_7-7)}[y_7(E_7 + n_7^2 y_7) + [y_7^2 (E_7 + n_7^2 y_7)^2 + 4 y_7 (E_7 - 7)(E_7 - n_7)]^{7/2}]$

Given $E_7 > n_7$ this lower limit to $X_7 / M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7 / M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$ in the appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7 < n_7$ $(ie., 2A_7 M_7^2 < X_7 - Y_7)$ then (A2.7)and (A2.8)show that $Q_7$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - (y_7 E_7 (7 + y_7)(E_7 - n_7)]^{7/2}]$$

And

$$E_7 > n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2 > 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.

If $A_7^2 < 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.+ $(Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$ and $X_7$ and $Y_7$ by $W_7^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSILIDATAN**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7\lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]\, e^{2T_7\lambda_7} + M_7 V_7\, e^{2T_7\lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i}\, Z_7$ *and* separate

(A3.1)into real and imaginary parts

$$Z_7^2 + (M_7 + V_7) Z_7 \sin(Z_7 T_7) - (U_7 M_7 + Y_7) + (M_7(U_7 + V_7) - X_7)$$

$$\mathrm{X}\cos(Z_7 T_7) - 2M_7 V_7 \cos^2(Z_7 T_7) + M_7 V_7 = 0$$

$$(M_7 + U_7) Z_7 - (M_7 + V_7) Z_7 \cos(Z_7 T_7) + (M_7(U_7 + V_7) - X)$$

$$X \sin\left(\left(Z_7 T_7\right) - M_7 V_7 \sin\left(Z_7 T_7\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_7$ and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - cos\ Z_7\ T_7$ and $A_7 = U_7 - V_7$ we obtain

$$4 M_7^2 V_7^2 W_7^4 + 2 M_7 V_7\ W_7^3 [(M_7 - V_7)^2 + A_7 (M_7 + V_7) + 2(A_7 M_7 - X_7)]$$

$$+ W_7^2 [M_7 - V_7)^2 (2 A_7\ M_7 + Y_7 - X_7) + (2A_7\ M_7 - X_7)^2 - A_7 (M_7 - V_7)(2A_7\ M_7 - X_7)$$

$$+ M_7 V_7 (A_7^2 + 4(X_7 + Y_7))]$$

$$+ W_7 [A_7^2 + 2X_7)(2 A_7\ M_7 - X_7) + 2A_7 Y_7 (M_7 + V_7) - M_7 A_7^3] + A_7^2 (X_7\ M_7 - Y_7) = 0$$

Being a quartic in $W_7$ (A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7 > 0$$

$$and$$

$$C_7 > (X_7 - Y_7)/2\ M_7$$

$recalling\ that\ A_7 = U_7$ and defining $C_7 = U_7 + V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0, U_7 V_7 = A_7^2 (X_7 + Y_7)$$

$$Q_7\left(2, U_7 V_7\right) = \left(C_7 + 2M_7^2\right)^2 (2C_7 M_7 + Y_7 - X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1\ in\ (A3.3)$

## ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

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LEN NUNNEYS

# Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price, University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss**

(1165) In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them.

(1166) Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors.

**Classical vs. relativistic black holes**

(1167) Something like a black hole exists within (eb) Newton's classical theory of gravity.

(1168) In that theory, an energy argument tells us that there is (=) an escape velocity vesc=2GM/R−−−−−−−√ from (e) the surface of any spherical object of (e) mass M and radius R

(1169) If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity.

(1170) Thus the condition for such an "unseeable" object is R<2GM/c2

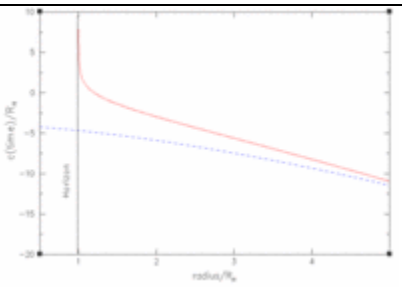
(1171) In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape.

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| (1172) This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity. |
| (1173) To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it. |
| (1174) The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c . |
| (1175) This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed. |
| (1176) An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity. |
| (1177) Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable. |

**Stationary black holes**

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| (1178) In the Newtonian theory, gravity is described by (e) the potential $\Phi$. |
| (1179) Inside a spherical object the form of (e) $\Phi(r)$ depends on (=) the interior structure, but in the vacuum outside (e) matter the potential $-GM/r$ depends only on (=) the interior mass. |
| (1180) Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object. |
| (1181) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward. |
| (1182) This horizon radius $RH=2GM/c2$ is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory. |
| (1183) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution. |
| (1184) Radius RHhere actually means that the area of the event horizon is (=) $4\pi R2H$. |
| (1185) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution. |
| (1186) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity. |
| (1187) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own. |
| (1188) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole. |
| (1189) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density. |
| (1190) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed. |
| (1191) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole. |

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| (1192) | In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter. |
| (1193) | Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum. |
| (1194) | In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior. |
| (1195) | The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object. |
| (1196) | Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum. |
| (1197) | These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.] |
| (1198) | All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time. |
| (1199) | But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?" |
| (1200) | The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point. |
| (1201) | One of these coordinates is (=) called the "coordinate time." |
| (1202) | Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate. |



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| (1203) | Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole |
| (1204) | At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole. |
| (1205) | Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations. |
| (1206) | For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree |

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| (1207) An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events. |
| (1208) The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving. |
| (1209) Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle. |
| (1210) The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return. |
| (1211) But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon. |
| (1212) Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future. |
| (1213) The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?" |

**Black hole parameters**

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| (1214) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin). |
| (1215) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun. |
| (1216) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits. |
| (1217) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form. |
| (1218) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds. |
| (1219) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2. |
| (1220) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed. |
| (1221) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer. |
| (1222) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2) |
| (1223) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials. |
| (1224) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable |
| (1225) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite. |
| (1226) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole. |

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| (1227) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars. |
| (1228) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity. |
| (1229) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity. |
| (1230) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite. |
| (1231) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity. |
| (1232) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum. |
| (1233) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship. |
| (1234) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= (4πG2/c4)[(M+M2−a2−−−−−−−√ )2+a2]. |
| (1235) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored. |

**Dynamical black holes**

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| (1236) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape. |
| (1237) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, **a horizon can be born**. |
| (1238) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid. |
| (1239) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward. |
| (1240) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations. |
| (1241) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes. |
| (1242) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3. |
| (1243) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/ |
| (1244) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers. |
| (1245) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission. |

(1246) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.

(1247) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.

(1248) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.

(1249) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.

(1250) The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature $6\times10-8$ (Mo/M) K.

(1251) Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.

**Astrophysical black holes**

(1252) Black holes in our Universe can be grouped as: primordial black holes, stellar-mass black holes, and supermassive black holes. The first is highly speculative; the second and third are broadly accepted.

(1253) Primordial black holes of all masses are postulated to have formed from (e) quantum fluctuations in the early Universe, but those with (e&eb) mass less than around 1013kg would have already evaporated due to Hawking radiation.

(1254) No significant observational evidence has yet been found of (e) the existence of these objects [MacGibbon and Carr (1991)].

(1255) Stellar-mass black holes, ranging from a few to a few tens of solar masses, are (=) normal but rare endpoints in (eb) the evolution of massive stars.

(1256) When a star exhausts its nuclear fuel and cools, it must collapse unless (e) it is supported by nonthermal forces.

(1257) It is known that such nonthermal forces cannot (e) resist gravitational compression for (e) masses greater than around 1.5 to 3Mo. (The uncertainty is due to (e) uncertainty in our knowledge of nuclear physics at high densities.)

(1258) There are many stars far more massive than this, but in their death throes massive stars expel (eb) a great deal of their mass in supernova explosions

(1259) Even stars initially as massive as 20-30Mo may blow off enough mass to (e) leave a remnant neutron star smaller than 1.5Mo.

(1260) Stars more massive than 20-30Mo may form (eb) a neutron star core, only to have it collapse due to (e) fallback of material from (e) the stellar mantle. Still more massive stars (above ~40Mo) may form (eb) a black hole directly in collapse, with or without (e&eb) a supernova explosion.

(1261) These masses are (=) rather uncertain at present, as our understanding of supernova explosion mechanisms is (=) still evolving [Woosley and Janka (2005), Muno (2006)].

(1262) Collapsing gas clouds larger than ~100Mo are expected to dissipate from (e) radiation pressure, which would prevent (e) more massive stars from forming, and consequently set (eb) an approximate upper limit for stellar mass black holes.

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| (1263) Observational evidence for stellar-mass black holes comes primarily from (e) X-ray astronomy. |
| (1264) A black hole in a close binary orbit can form (eb) an accretion disk (see Figure 1) of matter pulled off (er) a normal stellar companion. In its inspiral, disk material is heated by shearing, and becomes (=) a strong X-ray emitter. |
| (1265) If observations reveal a point-like X-ray emitter with (e&eb) a mass inferred from (e) orbital dynamics to be (=) above that possible for a neutron star, it becomes (=) a black hole candidate. |
| (1266) t present there are over a dozen such black hole candidates, including the first object to be identified as (=) a black hole, the X-ray source Cygnus X-1 (estimated mass 10±3Mo) [Casares (2006)]. |
| (1267) The other class of observationally-supported black holes is (=) that of supermassive black holes, ranging from hundreds to (e&eb) billions of solar masses. |
| (1268) Evidence for the existence of black holes at (eb) the upper end of this range is overwhelming. |
| (1269) By contrast, the existence of black holes with (e&*eb) mass roughly of order 102 to 104Mo, referred to as "intermediate-mass black holes," is speculative. |
| (1270) The supermassive holes may have begun as (=) primordial or stellar-mass black holes, but have grown through (e&eb) absorption of stars or gas or through mergers with (e&eb) other holes [Ferrarese and Ford (2005)]. |
| Figure 4: Observed orbits of stars near (eb) the dark compact object at the center of (e) our Milky Way galaxy: still frame from an animation. Animation created by Prof. Andrea Ghez and her research team at UCLA, from data sets obtained with the W. M. Keck Telescopes: http://www.astro.ucla.edu/~ghezgroup/gc/ |
| Evidence for supermassive black holes originally came from (e) quasars, small but intense radio sources seen at (eb) cosmological distances. |
| Their huge luminosities implied (eb) high mass, while (e) their rapid variability implied (eb) extremely small size. |
| More recently, measurements of Doppler shifts of stars and gas in the centers of galaxies have shown (eb) that compact objects of mass greater than 106Mo reside in (eb) the cores of most galaxies; the very small size of these central objects rules out any plausible alternative to the black hole explanation. |
| The black hole at the center of our own Milky Way galaxy has (e) a measured mass of ∼3×106Mo , based on (e) the orbits of stars near the radio source Sagittarius A* associated with (e&eb) the Galactic core [Ghez et al. (2005); see Fig. 4]. |
| (1276) Black holes are of particular interest for (e) gravitational waves, and vice versa. |
| (1277) For gravitationally-bound systems, the typical maximum gravitational wave amplitude (dimensionless strain) is (=) h∼(2GM/c2)2/(RD)∼R2H/(RD) , where D is the distance to the system, M and R are the mass and size of (e) the system, and RH is the horizon radius of (e) a black hole of the same mass. |
| (1278) From this it is clear that strongest gravitational waves will involve (e&eb) systems at or near black hole compactness: in particular, systems containing (er) black holes in close proximity to one another [Flanagan and Hughes (2005)]. |
| (1279) Conversely, a system containing only black holes would not be expected to radiate (eb) any sort of radiation other than gravitational waves, and in any system containing (e) black holes, gravitational waves provide (eb) the most accurate probes of conditions deep in (eb) their gravitational wells. |
| (1280) In particular, gravitational wave observations could definitively test (e&eb) whether a compact mass is truly a black hole [Cutler and Thorne (2002)]. |

(1281) The strongest intrinsic source of gravitational waves in the present Universe is (=) the collision of two comparable-massed black holes.

(1282) A sufficiently tight black hole binary system will lose (e) energy through emission of (e&eb) gravitational waves, causing (eb) the orbit to circularize and then slowly shrink over time.

(1283) Figure 5: Gravitational waves from (e) a black hole binary inspiral: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/

(1284) When the orbital radius approaches (eb) the Schwarzschild radius of the system, complex nonlinear dynamics come into play (e&eb) and the final stages of (e) the inspiral and merger can only be modeled through (e&eb) numerical simulations.

(1285) These simulations are quite challenging, but recent breakthroughs [Pretorius (2005)(See Fig. 5), Campanelli et al.(2006), Baker et al.(2006)] have led to (eb) the first complete inspiral-merger gravitational waveforms.

(1286) After the black holes have coalesced, they form (eb) a single highly-distorted black hole which quickly settles down to (e) a quiescent Kerr state through (e&eb) mission of quasinormal "ringdown" gravitational waves [Pretorius (2005), Campanelli et al. (2006), Baker et al. (2006)].

(1287) Another gravitational emission mechanism for supermassive black holes would come from (e) their capture of compact stellar objects, particularly stellar-mass black holes.

(1288) Unlike comparable-mass mergers, these captures would be on (eb) highly eccentric orbits that would not circularize before merger, resulting in (eb) intricate rapidly precessing orbits in (eb) the final months or years of inspiral.

**(1289)** The consequent gravitational-wave signal would allow (eb) for exquisitely precise measurements of the parameters of (e) the supermassive black hole, as well as tests of any deviation from (e) the predictions of general relativity [Cutler and Thorne (2002)].

**(1290)** For stellar-mass black holes (tens of Mo), the final stages of inspiral and merger occur at (eb) frequencies of hundreds to thousands of Hz, a regime targeted by (e) ground-based gravitational-wave observatories.

**(1291)** Systems of supermassive black holes that are (=) millions of Mo emit at frequencies of tens of (e) mHz, which will be detectable by (e) proposed space-based (e) gravitational-wave detectors.

**(1292)** Black holes in the billions of Mo, or a stochastic background of signals from (e) more moderate supermassive black holes at (eb) earlier stages of their inspiral, might emit (eb) waves at nHz frequencies that would be measurable through (e&eb) years-long correlated timing measurements of radio pulsars [Jenet et al. (2005)].

**An equation to quantify the origins of life on other planets July 5, 2016 by Bob Yirka report**

(1293) (Phys.org)—A pair of researchers, one with the Columbia Astrobiology Center in New York, the other with the University of Glasgow in the U.K. has come up with a mathematical equation that when (e) solved is meant to offer (eb) a means for estimating how often life begins on other planets. In their paper published in Proceedings of the National Academy of Sciences, Caleb Scharf and Leroy Cronin describe their equation, how they came up with it and why they believe it might become more useful as scientists learn more about the true nature of other planets and solar systems.

(1294) Back in the 1960's, an astronomer named Frank Drake came up with a formula for (e) estimating the likely number of (e) alien civilizations that might be capable of transmitting radio signals in such a way as to be (=) recognizable by receivers here on earth.

(1295) That equation, named quite naturally the **Drake Equation** has been the unofficial standard-bearer for half a century, despite it having no solution because (e) some of the parameters are still unknown.

(1296) The same holds true for the new equation developed by (e) Scharf and Cronin, but the results given should the unknowns ever be discovered would offer a different sort of answer than Drake was looking for.

(1297) To build their formula, the duo came up with several parameters to solve for (e) **abiogenesis**, which is the likelihood of an event occurring that leads to (eb) life beginning, they included (e) the number of possible building blocks, the mean number of such blocks per possible organism, the availability of building blocks that might exist during a given time period, expressed as a fraction and the probability that the existence of the building blocks would actually lead to (eb) life starting in a given unit of time. It looks like this: **Nabiogenesis (t)= Nb · 1/no · fc · Pa · t**

(1298) In essence the formula is meant to suggest (eb) that the probability of life beginning on a given planet is very likely connected to (e&eb) whether there are building blocks available on a given planet, and how much of them there might be.

(1299) As more is learned by space researchers it is hoped that the formula could help narrow down (e) search targets by offering (eb) a statistical probability of success for (e) a given planet.

(1300) The researchers note that the building blocks available do not necessarily have to be (=) of the type that led to (eb) life beginning on our planet.

**(1301)** They also suggest that implied in the formula is (=) the likelihood that life beginning events are (=) more likely to occur in solar systems where (e) there are multiple planets, which allow (eb) for an opportunity to share materials that could lead to (eb)building blocks **More information: Quantifying the origins of life on a planetary scale, Proceedings of the National Academy of Sciences, www.pnas.org/cgi/doi/10.1073/pnas.1523233113**

**Abstract**

(1302) A simple, heuristic formula with (e&eb) parallels to the Drake Equation is introduced to help focus discussion on open questions for (e) the origins of life in a planetary context.

(1303) This approach indicates (eb) a number of areas where (e) quantitative progress can be made on parameter estimation for (e) determining origins of life probabilities, based on (e) constraints from Bayesian approaches.

(1304) Authors discuss a variety of "microscale" factors and their role in determining (eb) "macroscale" abiogenesis probabilities on (e&eb) suitable planets.

(1305) Authors also propose that impact ejecta exchange between (e&eb) planets with parallel chemistries and (e&eb) chemical evolution could in principle amplify (eb) the development of (e) molecular complexity and (e&eb) abiogenesis probabilities.

(1306) This amplification could be very significant, and both bias conclusions about abiogenesis probabilities based on (e) the Earth and provide (eb) a major source of variance in the probability of life arising in (eb) planetary systems.

**(1307)** Authors use heuristic formula to suggest (eb) a number of observational routes for improving constraints on (e) origins of life probabilities. **Read more at: http://phys.org/news/2016-07-equation-quantify-life-planets.html#jCp**

**Brian Koberlein Shared publicly - Yesterday 8:00 PM the Illusion of Truth**

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| (1308) | Physical theories are often presented as (=) a description of what's really going on. |
| (1309) | Forces act on (e&eb) a baseball, causing (eb) baseball to fall. |
| (1310) | Atoms collide and fuse in (eb) the heart of a star, releasing heat and energy. Science is true, as is often said, and our scientific theories encapsulate this truth. But this isn't entirely true. |
| (1311) | For example, quantum mechanics is a strange, sometimes confusing theory. Objects can be (=) particles and waves. |
| (1312) | They do strange things like tunnel through (e&eb) barriers, and can even appear and disappear (e&eb). |
| (1313) | Given all this strangeness, what's really going on? What's the true nature of quantum reality? It depends on (=) which approach you want to take. |
| (1314) | In one view quantum objects are described by (e) a wave-like probability function. |
| (1315) | When you interact with (e&eb) or "observe" (e) the object this "wavefunction" collapses into (eb) a definite state. |
| (1316) | That's view commonly presented, but a quantum system can also be described by (e) its transitions between energy states. |
| (1317) | Since this uses a mathematical method involving (e&eb) matrices, it's known as (=) matrix mechanics. |
| (1318) | Both wavefunctions and matrix mechanics give (eb) the same results, but their view of what really goes on is very different. |
| (1319) | Then there's the path integral method |
| (1320) | . Rather than a wavefunction or matrix transitions, path integrals imagine quantum objects can take almost any path between (e&eb) two states. |
| (1321) | By summing all the possible paths you can derive (eb) the odds that it will occur. |
| (1322) | So which one is true? Are quantum objects distributed waves of (e) probability? |
| (1323) | Are they simply transitions between (e&eb) energy states? |
| (1324) | Do they take an infinite number of (e) paths between states? |
| (1325) | Each of these models makes (eb) the same predictions, so one version is no more "true" than the others. |
| (1326) | What happens in practice is that we'll use whatever method is useful at the time. They are equivalent models, so the best model for the job is the one we'll use. The only reason the wavefunction view is so common is that it's (=) the version usually taught to introductory students. You might think this uncertainty of truth is due to (e) the behavior of quantum physics itself. It's so strange and counterintuitive that we can't (e) wrap our puny brains around what's really going on. But the same thing occurs in lots of other fields. Even something as straight forward as basic Newtonian physics |
| (1327) | Toss a baseball in the air and (eb) Earth's gravitational force pulls it down. |
| (1328) | The force of gravity is a simple truth, right? While we often describe classical motion in terms of (e&eb) forces and acceleration, we can also describe it in terms of (e&eb) energy and momentum. |
| (1329) | In the Lagrangian and Hamiltonian approach, the path of a baseball is (=) the optimized path among possibilities. |
| (1330) | In this view a baseball's path is (=) the extrema of an energy equation, and force can be derived as a necessary consequence of (e) this. |
| (1331) | In the relativistic view the baseball follows (e&eb) a geodesic, which is (=) the minimal path through (e&eb) space and time. |

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| (1332) So is a baseball's motion due to (e) a gravitational force, an energy extrema, or a spacetime geodesic? |
| (1333) As with quantum theory, different approaches yield (eb) the same result. They are mathematically equivalent, so we can use whatever method is most useful at the time. |
| (1334) At its core, science is less about truth and more about models. The metaphysics underlying a model is useful only as far as it allows us (eb) to make better predictions, generate (eb) new ideas, or bring models together as (=) a cohesive whole |
| (1335) This is why we have no problem using (e) classical gravity to calculate the path of (e) a spacecraft through (e&eb) the solar system, while (e) using special relativity to account for (e) the Doppler shift of the spacecraft's radio signals. |
| (1336) It's why we can use (e) quantum physics to study atoms in the morning, and general relativity to (e) study black holes in the afternoon. |
| (1337) In regimes where models conflict with each other it isn't a failure of truth, but instead shows an opportunity to develop a better model. |
| (1338) It could be that with each better model we move closer to (e&eb) the truth about reality. |
| (1339) The truth is out there, and science strives to move towards (e&eb) that truth. |
| (1340) It's a common view, and certainly the search for truth has driven (eb) many scientists to develop better and better models. |
| (1341) But the real power of science is the recognition that what we have are models. Our models can be powerful, but they are always a bit tentative. There's always a chance that they might just be an illusion of truth. |

**Synopsis: A Single-Photon Cheshire Cat July 1, 2016 Researchers detected the polarization of a photon separate from the photon itself, just as the grin of Lewis Carroll's Cheshire cat can appear apart from the cat's body Synopsis figure APS/Carin Cain**

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| (1342) In a quantum Cheshire cat, one property of a particle is present detached from (e) the particle itself, just as (=) the grin of the feline in Alice in Wonderland can exist in the absence of (e) the whole animal. |
| (1343) Now a research team has successfully separated (e&eb) a single photon from (e) its polarization. |
| (1344) The effect was previously demonstrated (eb) with neutrons and also with light in a classical scenario, but the latest experiment uses single photons. |
| (1345) To produce the Cheshire cat, **Maximilian Schlosshauer and his colleagues** at the University of Portland in Oregon used an interferometer where (e) light hits a beam splitter, goes in two, opposite directions around a circuit (but along slightly offset routes), and then recombines at (eb) the entrance. |
| (1346) The incoming polarization of light was horizontal, but this was switched to (e&eb) vertical for the counterclockwise loop. |
| (1347) Each path included a device to adjust (e&eb) the polarization further and another to (e&eb) scatter away a small number of photons of (e) a designated polarization, the equivalent of detecting their presence. |
| (1348) The team fired single photons into (e&eb) the interferometer so that only one photon was present at (eb) a time, traversing both paths simultaneously. |
| (1349) They placed a horizontal polarizer in front of the output photon detector and gradually varied (e&eb) the difference between the two path lengths. |

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| (1350) This variation generated (eb) a sine-wave interference pattern from (e) the exiting photons only when the polarization was tweaked in (eb) the counterclockwise path. |
| (1351) At the same time, the average number of photons counted decreased only when (e) light was scattered from the clockwise path. |
| (1352) Simultaneously demonstrating these different effects for two routes of a single photon constitutes (eb) evidence of the Cheshire cat. |
| **(1353)** The effect might someday improve precision experiments if (e) an unwanted property could be moved away from the property being measured. **This research is published in Physical Review A. – David Ehrenstein David Ehrenstein is the Focus Editor for Physics.** |

**Observation of the quantum paradox of separation of a single photon from one of its properties James M. Ashby, Peter D. Schwarz, and Maximilian Schlosshauer Phys. Rev. A 94, 012102 – Published 1 July 2016**

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| (1354) Authors report an experimental realization of the quantum paradox of (e) the separation of a single photon from (e) one of its properties (the so-called "quantum Cheshire cat"). |
| (1355) Authors use a modified Sagnac interferometer with (e&eb) displaced paths to produce (eb) appropriately pre- and postselected states of (e) heralded single photons. |
| (1356) Weak measurements of photon presence and (e&eb) circular polarization are performed in each arm of (e) the interferometer by introducing (eb) weak absorbers and small polarization rotations and analyzing changes in (eb) the postselected signal. |
| (1357) The absorber is found to have (e) an appreciable effect only in one arm of the interferometer, while (e) the polarization rotation significantly affects (e&eb) the signal only when (e) performed in the other arm. |
| (1358) Authors carry out both sequential and simultaneous weak measurements and find (eb) good agreement between measured and predicted weak values. |
| (1359) In the language of Aharonov et al. and in the sense of the ensemble averages described by (e) weak values, the experiment establishes (eb) the separation of a particle from one its properties during the passage through (e&eb) the interferometer. |

**Viewpoint: Relativity Gets Thorough Vetting from LIGO Frans Pretorius, Department of Physics, Princeton University, Princeton, NJ 08544, USA May 31, 2016• Physics 9, 52 The gravitational-wave signal observed by the LIGO detectors shows no deviation from what general relativity predicts. Figure caption expand figure Frans Pretorius, APS/Carin Cain**

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| (1360) Figure 1: The signal from one of the LIGO detectors in Hanford, Washington, is shown with two representations of (e) the best-fit numerical relativity (NR) waveform. |
| (1361) The filtered NR waveform illustrates how the raw waveform is perceived by (e) the detector, |
| (1362) On September 14, 2015, the LIGO interferometers made the first direct measurement of (e) gravitational waves [1]. |
| (1363) The historic discovery came nearly 100 years after Einstein showed that his general relativity theory predicts (eb) the existence of gravitational waves and over 50 years since (e) the first quest to observe the Universe in gravitational waves [2]. |
| (1364) The observed event, labeled GW150914, was (=) the collision of two black holes, one weighing 29 solar masses, the other 36 solar masses. |

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| (1365) This binary coalesced into (e&eb) a remnant black hole weighing 62 solar masses and rotating (e&eb) at 67% the maximum rate allowed (eb) for a black hole. |
| (1366) The most statistically significant portion of (e) the signal came from (e) the last several gravitational-wave cycles of the collision, where (e) the most nonlinear, strong-field dynamics of the theory manifest. |
| (1367) One couldn't have asked for a better first event to learn about (e&eb) this regime of general relativity. |
| (1368) The first thing that GW150914 can teach us is how closely nature follows (eb) Einstein's prescriptions for (e&eb) colliding black holes. |
| (1369) Researchers from the LIGO team and the Virgo team have performed a series of tests to see if (e) any violations of (e) general relativity might be lurking in the data of GW150914 [3]. |
| (1370) They first determined (eb) the theoretical waveform, or "template," from (e) general relativity simulations that best fit the measured signal (see Fig. 1). |
| (1371) Their main finding is that the residual signal—obtained by (e) subtracting this template from (e) the data—is consistent with (e&eb) noise. |
| (1372) This is, in a sense, the "everything" test, as (e) the template folds in (e&eb) all that general relativity predicts about the merger: the decay of the last few orbits of (e) the two black holes, their violent collision (e&eb) and astonishingly rapid relaxation to (e) a quiescent rotating black hole, the propagation of (e) gravitational waves through (e&eb) the cosmos, and finally the transformation of (e) these waves into(e&eb) a chirp in the LIGO detectors. |
| (1373) Given the large signal-to-noise ratio of 24, the close agreement between the template and (=) the data puts limits on (e) alternative models. |
| (1374) Specifically, the predictions of these models can only disagree with (e&eb) the general relativity prediction for GW150914 by, at most, 4%. |
| (1375) In addition to this comprehensive test, the LIGO/Virgo teams also performed other checks that pick apart the signal, in order to address particular aspects of general relativity and black hole physics. In one such test, they compared the early and (e&eb) late parts of the waveform. |
| (1376) From the waves generated during (e&eb) the initial inspiral, one can infer (eb) the masses and spins of the binary. |
| (1377) Given these numbers, general relativity allows (eb) the mass and spin of (e) the remnant black hole to be calculated. |
| (1378) The researchers compared this result to (e&eb) an independent measurement of (e) the mass and spin obtained from (e) the late-time ringdown, and to within (eb) uncertainties, the two estimates were (=) consistent. |
| (1379) Another test isolated coefficients in the so-called Post Newtonian (PN) series describing (eb) the inspiral. |
| (1380) The PN series is (=) an expansion of the gravitational-wave signal in powers of (e) the velocity vv of the sources relative to (e&eb) the speed of light cc. |
| (1381) Observations of binary pulsars, which show (eb) a decrease in their orbital period due to (e) gravitational-wave emission, have constrained (e) PN terms, especially at low order. |
| (1382) For higher-order terms associated with (e&eb) strong-field emission, the LIGO/Virgo teams were able to improve (eb) the limits by many orders of magnitude compared to (e&eb) the binary pulsar constraints. |

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| (1383) At first glance it might seem these new results simply add to (e) the litany of previous tests that general relativity passed with flying colors. |
| (1384) However, GW150914 is vastly more significant than that. Other tests have shown that general relativity is valid in (eb)the weak-field regime, but GW150914 gives (eb) the first direct observational evidence that general relativity is the correct theory of nature governing (e&eb) the extreme spacetime dynamics around the event horizon of (e) a black hole. |
| (1385) One could argue that binary pulsar observations already told us that gravitational waves exist [4], but GW150914 is more luminous in gravitational waves than pulsar binaries by (e) at least 24 orders of magnitude. |
| (1386) It's true that the evolution of the Universe as a whole is driven by (e) strong gravity, but it would be difficult to argue that precision cosmological measurements can constitute (eb) a test of general relativity when (e) one of the dominant inputs, dark energy, is still a mystery. |
| (1387) As far as black holes are concerned, astronomers have (e) observed ultra compact dark objects that are consistent with (e&eb) black holes. |
| (1388) The most striking data comes from (e) stellar orbits about the 4 million solar mass objects in (eb) the center of our Galaxy [5], but present observations cannot constrain it to be (=) a black hole as described by general relativity. |
| (1389) GW150914 has given us the first morsels of quantifiable evidence that the rich physics of black holes is as predicted by general relativity. For example, the LIGO observations are consistent with (e&eb) so-called cosmic censorship [6], which says that gravitational singularities must remain hidden behind an (eb) event horizon. |
| (1390) If a "naked" singularity were produced in (eb) the energetic collision of GW150914, the data would likely have shown (eb) some sort of anomaly, but in exactly what form, we don't know. |
| (1391) Another example involves the "no-hair" theorem [7], which says that the only properties of a black hole are (=) its mass, spin, and charge (though astrophysical black holes are expected to have negligible charge). |
| (1392) By consequence, these properties should determine (eb) the frequency and decay constants of (e) all the infinitely many ringdown modes of (e) a black hole. |
| (1393) With GW150914, only the least-damped ringdown mode is (=) visible above the noise, but when combined with (e&eb) the remnant mass and spin inferred from (e) the inspiral, this is enough to count as (=) the first piece of evidence for a black hole's baldness. |
| (1394) GW150914 should also be a test for (e) alternative theories of gravity or exotic alternatives to (e) black holes within general relativity, such as boson stars [8]. |
| (1395) Some of these alternative models are concerned with (e&eb) explaining dark energy or dark matter and do not (e) therefore, predict any deviation from relativity at (eb) the scale of GW150914. |
| (1396) However, other models do in fact deviate from (e) conventional relativity in their predictions of (e) the gravitational-wave signal from (e) a binary merger. |
| (1397) The largest deviations should presumably occur during (e&eb) the last several gravitational-wave cycles, where (e) the signal is strong, and deviations should be (=) easy to detect or constrain. |
| (1398) Embarrassingly though, the majority of alternatives that claim deviations can't (e) really say much about this event, as the two-body problem hasn't been solved in these theories. |
| (1399) These models can only reproduce (eb) a fraction of the full gravitational-wave signal, (eb) so the strong constraints one might have hoped for (e) cannot be leveled. |

(1400) Exotic alternatives within (eb) general relativity are in a similar boat, or worse for (e) objects like gravastars, which have (e) no well-defined theoretical framework to say what happens when (e) two such objects collide.

(1401) Regardless, what is clear from the data is that any non-black-hole solution must imagine a form of (e) exotic matter with a viscosity larger than (e) any known material to explain (eb) the rapid damping of oscillations observed during the ringdown.

(1402) The inferred event rate for heavy black hole mergers like GW150194 is high, and LIGO has another factor of 3 improvement in sensitivity slated before it reaches (eb) design sensitivity.

(1403) Within 5 years then, we can expect dozens of detected events, with the loudest having a signal-to-noise ratio upwards of (e) 100.

(1404) This will allow (eb) for exquisite insight into the nature of extreme gravity beyond (e) what GW150194 has already gifted us with. **This research is published in Physical Review Letters**.

**Tests of General Relativity with GW150914 B. P. Abbott et al. (LIGO Scientific and Virgo Collaborations) Phys. Rev. Lett 116, 221101 – Published 31 May 2016**

(1405) The LIGO detection of GW150914 provides (eb) an unprecedented opportunity to study the two-body motion of (e) a compact-object binary in (eb) the large-velocity, highly nonlinear regime, and to witness the final merger of (e) the binary and the excitation of uniquely relativistic modes of (e) the gravitational field.

(1406) Authors carry out several investigations to determine whether GW150914 is consistent with (e&eb) a binary black-hole merger in general relativity.

(1407) They find that the final remnant's mass and spin, as determined from (e) the low-frequency (inspiral) and (e&eb) high-frequency (postinspiral) phases of the signal, are mutually consistent with (e&eb) the binary black-hole solution in general relativity.

(1408) Furthermore, the data following the peak of GW150914 are consistent with (e&eb) the least-damped quasinormal mode inferred from (e) the mass and spin of the remnant black hole.

(1409) By using waveform models that allow (eb) for parametrized general-relativity violations during (e&eb) the inspiral and merger phases, authors perform quantitative tests on (e&eb) the gravitational-wave phase in the dynamical regime and they determine the first empirical bounds on (e&eb) several high-order post-Newtonian coefficients.

(1410) Authors constrain (e) the graviton Compton wavelength, assuming (eb) that gravitons are dispersed in (eb) vacuum in the same way as particles with mass, obtaining a 90%-confidence lower bound of 1013 km.

(1411) In conclusion, within our statistical uncertainties, we find no evidence for violations of (e) general relativity in the genuinely strong-field regime of gravity. **Received 26 March 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.221101 © 2016 American Physical Society**

**Focus: Background Noise of Gravitational Waves March 31, 2016• Physics 9, 33**

(1412) Assuming LIGO's recently detected black hole merger is not unusual; researchers revised upward their estimate of (e) the strength of the background noise coming from (e) distant mergers across the Universe. Figure caption expand figure Caltech/MIT/LIGO Lab (deleted due to spatial constraints)

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| (1413) Catching some waves The LIGO-Livingston interferometer consists of (e) two 4-km-long arms where (e) laser light bounces back and forth. Interference between the lasers from the two arms can signal tiny changes in (e&eb) the lengths of the arms resulting from gravi... |
| (1414) The recent detection of gravitational waves emitted from (e) a pair of merging black holes implies (eb) the existence of a large population of similar pairs, which should be generating (eb) a background of waves too faint to be detected with (e&eb) present technology. |
| (1415) Using this information, researchers now estimate that this background could be (=) roughly 10 times stronger than previously expected and might be found by (e) sensitive instruments in only a few years. |
| (1416) The background signal could help astrophysicists better understand how black hole binaries form and also pave the way to detecting the waves emitted by (e) the expanding Universe shortly after the big bang. |
| (1417) Gravitational waves were detected for the first time on 14 September 2015 by the Laser Interferometer Gravitational-Wave Observatory (LIGO). |
| (1418) By reflecting laser light between mirrors set 4 km apart, the team was able to detect (e&eb) tiny fluctuations in the distance between the mirrors as (=) the waves passed. |
| (1419) The waves were triggered by (e) an event, called GW150914; involving (e&eb) the merger of a pair of black holes having (e) masses of about 29 and 36 times that of the Sun. |
| (1420) Researchers expect that similar pairs of black holes, as well as neutron stars, should be merging constantly, creating (eb) gravitational waves that should bombard (e&eb) the Earth from all directions and make the equivalent of (=) a faint popping noise in a detector. |
| (1421) Previous estimates of the likely strength of this background had to make assumptions about (e&eb) how many pairs there might be and about the distribution of their masses. |
| (1422) These estimates remained speculative, as no actual binaries were known. |
| (1423) “What the recent observation has changed,” says Nelson Christensen of Carleton College in Minnesota and the LIGO Scientific Collaboration (LSC), “is that we now know there really is (=) a population of black hole binaries. We have a small bit of information about it.” |
| (1424) The new analysis—carried out by LSC researchers working with (e&eb) the European Virgo collaboration—estimates (eb) the energy density in the gravitational wave background in (eb) the frequency range below 100 Hz, where (e) the LIGO detector is most sensitive. |
| (1425) To make this estimate, they assumed (eb) a population of black hole binaries of various sizes, having (e) a range of distances from Earth, and created through (e&eb) a variety of scenarios. |
| (1426) They then calculated the amount of gravitational radiation that would result from (e) the merging of these objects. |
| (1427) The team then adjusted this estimate by using (e) details of the GW150914 event—especially the masses of the black holes involved—to make (eb) a most likely guess at the actual distribution of black hole binaries. |
| (1428) The idea is that the observed merger is (=) most likely representative of (e) the general population, rather than being extremely unusual. |
| (1429) If the black hole binary population is represented as (=) a bell curve, or Gaussian, the event should lie within (eb) the main bulge, not in the tails. |
| (1430) The observed mass of GW150914 suggests (eb) that the main population of black hole binaries could be (=) roughly 20 times more massive than expected, which would generate (eb) a background roughly 10 times “louder” than (e) previous estimates. |

(1431) However, Christensen emphasizes that the result has (e) lots of uncertainty, as it's based on (e) only one event.

(1432) "This estimate seems very plausible to me," says astrophysicist Maura McLaughlin of the University of West Virginia in Morgantown. She is a member of the NANOGRAV collaboration, which tries to detect gravitational waves at much lower frequencies than LIGO can observe. "I can't see any holes in their arguments. Unless LIGO got really lucky, then there must be many more of these systems than previously thought." This level of gravitational background implies (eb) that the LIGO and Virgo detectors might be sensitive enough to see it in 2–5 years, sooner than physicists previously expected, says Christensen, which has led to much anticipation. "Everyone is super excited," he says.

(1433) Detecting the background would provide (eb) a confirmation of the presence of gravitational waves carrying (e&eb) energy throughout the Universe. It would also be useful, the researchers suggest, in gaining (e) further understanding of precisely how black holes form.

(1434) And subtracting this background might someday allow (eb) detection of the fainter waves coming from (e) the expansion of the Universe moments after the big bang. This research is published in Physical Review Letters. –Mark Buchanan Mark Buchanan is a freelance science writer in Normandy, France.

**GW150914: Implications for the Stochastic Gravitational-Wave Background from Binary Black Holes B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 131102 – Published 31 March 2016**

(1435) The LIGO detection of the gravitational wave transient GW150914, from (e) the inspiral and merger of (e) two black holes with masses $\gtrsim 30 M_{\odot}$, suggests (eb) a population of binary black holes with (e&eb) relatively high mass.

(1436) This observation implies (eb) that the stochastic gravitational-wave background from binary black holes, created from (e) the incoherent superposition of (e) all the merging binaries in the Universe, could be (=) higher than previously expected.

(1437) Using the properties of GW150914, authors estimate the energy density of (e) such a background from (e) binary black holes.

(1438) In the most sensitive part of the Advanced LIGO and Advanced Virgo band for (e) stochastic backgrounds (near 25 Hz), authors predict $\Omega_{GW}(f=25\ \mathrm{Hz}) = 1.1^{+2.7}_{-0.9} \times 10^{-9}$ with 90% confidence.

(1439) This prediction is robustly demonstrated for (e) a variety of formation scenarios with (e&eb) different parameters.

(1440) The differences between models are (=) small compared to the statistical uncertainty arising from (e) the currently poorly constrained (e) local coalescence rate.

(1441) Authors conclude that this background is potentially measurable by (e) the Advanced LIGO and Advanced Virgo detectors operating at their projected final sensitivity. DOI:http://dx.doi.org/10.1103/PhysRevLett.116.131102 © 2016 American Physical Society

**Study opens new realms of light-matter interaction July 15, 2016 by David L. Chandler. Some sentences might have been omitted due to spatial constraints. Please bear with me.**

(1442) Emission spectra are a widely used method for (e) identifying chemical compounds; the bright lines reveal (eb) the different frequencies of light that can be emitted by (e) an atom.

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| (1443) Here, a normal emission spectrum for (e) an atom in a high-energy state (top) is compared to (e&eb) the emission from the same atom placed just a few nanometers (billionths of a meter) away from (e) graphene that has been doped with (e&eb) charge carriers (bottom). |
| (1444) For each energy-level transition, an orange line (or purple cloud) appears if (e) that transition is estimated to be faster than (e) one per microsecond — making it frequent enough to be (=) observed. **Credit: Massachusetts Institute of Technology** |
| (1445) A new MIT study could open up new areas of technology based on (e) types of light emission that had been thought to be (=) "forbidden," or at least so unlikely as to be (=) practically unattainable. |
| (1446) The new approach, the researchers say, could cause (eb) certain kinds of interactions between light and (e&eb) matter, which would normally take billions of years to happen, to take place instead (e&eb) within billionths of a second, under (e&eb) certain special conditions. The findings, based on a theoretical analysis, are reported today in the journal Science in a paper by MIT doctoral student Nicholas Rivera, Department of Physics Professor Marin Soljačić, Francis Wright Davis Professor of Physics John Joannopoulos, and postdocs Ido Kaminer and Bo Zhen. |
| (1447) Interactions between light and matter, described by (e) the laws of quantum electrodynamics, are the basis of a (e) wide range of technologies, including (e) lasers, LEDs, and atomic clocks. |
| (1448) But from a theoretical standpoint, "Most light-matter interaction processes are 'forbidden' by (e) electronic selection rules, which limits (e) the number of transitions between energy levels we have access to," Soljačić explains. |
| (1449) For example, spectrograms, which are used to (e) analyze the elemental composition of materials, show (eb) a few bright lines against (e&eb) a mostly dark background. |
| (1450) The bright lines represent (eb) the specific "allowed" energy level transitions in (eb) the atoms of that element that can be accompanied by (e) the release of a photon (a particle of light). |
| (1451) In the dark regions, which make up most of the spectrum, emission at those energy levels is (=) "forbidden." |
| (1452) With this new study, Kaminer says, "we demonstrate theoretically that these constraints can be lifted" using (e) confined waves within (e&eb) atomically thin, 2-D materials. |
| (1453) "We show that some of the transitions which normally take the age of the universe to happen could be made to happen within (eb) nanoseconds. |
| (1454) Because of this, many of the dark regions of a spectrogram become (=) bright once an atom is (=) placed near a 2-D material." |
| (1455) Electrons in an atom have (e) discrete energy levels, and when they hop from one level to another they give off (eb) a photon of light, a process called spontaneous emission. |
| (1456) But the atom itself is much smaller than the wavelength of the light that gets emitted—about 1/1,000 to 1/10,000 as big—substantially impairing (e&eb) the interactions between the two. |
| (1457) The trick is, in effect, to "shrink" (e) the light so it better matches the scale of the atom, as (=) the researchers show in their study. |
| (1458) The key to enabling a whole range of interactions, specifically transitions in (eb) atomic states that relate to absorbing or emitting light, is (=) the use of a two-dimensional material called graphene, in which light can interact with (e&eb) matter in the form of plasmons, a type of electromagnetic oscillation in the material. |
| (1459) These plasmons, which resemble photons but have wavelengths hundreds of times shorter, are (=) very narrowly confined in the graphene, in a way that makes some kinds of interactions with (e&eb) that matter many orders of magnitude more likely than (e) they would be in ordinary materials. |

(1460) This enables a variety of phenomena normally considered unattainable, such as (e) the simultaneous emission of multiple plasmons, or two-step light-emitting transitions between (e&eb) energy levels, the team says.

(1461) This method can enable the simultaneous emission of (e) two photons that are "entangled," meaning (eb) they share the same quantum state even when (e) separated.

(1462) Such generation of entangled photons is (=) an important element in quantum devices, such as those that might be used for cryptography.

(1463) Making use of these forbidden transitions could open up (eb) the ability to tailor the optical properties of materials in ways that had not been thought possible, Rivera says.

(1464) "By altering these rules" about the relationship between light and (e&eb) matter, "it can open (eb) new doors to reshaping (e&eb) the optical properties of materials."

(1465) Kaminer predicts that this work "will serve as (=) a founding piece for (e) the next generation of studies on light-matter interactions" and could lead to (eb) "further theoretical and experimental advances in many fields which rely on (e) light-matter interactions, including (e) atomic, molecular and optical physics, photonics, chemistry, optoelectronics, and many others."

(1466) Beyond its scientific implications, he says, "this study has possible applications across multiple disciplines, since (e) in principle it has potential to enable the full use of (e) the periodic table for optical applications."

(1467) This could potentially lead to (eb) applications in spectroscopy and sensing devices, ultrathin solar cells, and new kinds of materials to absorb solar energy, organic LEDs with higher efficiencies, and photon sources for (e) possible quantum computing devices."From the standpoint of fundamental science, this work lays the groundwork for a subfield that just a few years ago was difficult to imagine and until now was largely unexplored," Soljačić says.

(1468) "Two-dimensional materials confine (e) fields to a surface and motion to a plane, making possible (eb) many effects that are orders-of-magnitude too weak to (e) appear in a bulk volume," says Jason Fleischer, an associate professor of electrical engineering at Princeton University, who was not involved in this research.

(1469) This work, he says, "systematically explores how 2-D materials improve light-matter interactions, laying a theoretical foundation for faster electronic transitions, enhanced sensing, and better emission, including the compact generation of broadband and quantum light." More information: N. Rivera et al: Shrinking light to allow (eb) forbidden transitions on the atomic scale, Science (2016). DOI: 10.1126/science.aaf6308 Journal reference: Science search and more info website: Provided by: Massachusetts Institute of Technology search and more info website. This story is republished courtesy of MIT News (web.mit.edu/news office/), a popular site that covers news about MIT research, innovation and teaching.

**Viewpoint: Peering Back to the Beginning of Time Lawrence M. Krauss, Department of Physics and School of Earth and Space Exploration, Arizona State University, Tempe, AZ 85287, USA and Mount Stromlo Observatory, Research School of Astronomy and Astrophysics, Australian National University, Weston, ACT 2611, Australia June 19, 2014• Physics 7, 64**

(1470) The BICEP2 collaboration reports the detection of BB-mode polarization of (e) the cosmic microwave background—a signal that might originate from (e) gravitational waves created by (e) inflation during the very earliest moments in the evolution of our Universe. Figure caption expand figure P. A. R. Ade et al. [1]; Image on homepage: BICEP2 Collaboration (deleted due to restrictions on space)

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| (1471) Figure 1: (Top) BB-mode polarization pattern observed by (e) the BICEP2 telescope [1]. |
| (1472) The twisting pattern in the polarization of (e) the cosmic microwave background (CMB) might originate from (e) gravitational waves created by inflation in the very first moments of the universe. |
| (1473) Inflation—the hypothesis that the Universe underwent (e&eb) a phase of superluminal expansion in a brief period following the big bang—has (e) the potential of explaining, from first principles, why the Universe has (e) the structure we see today. |
| (1474) It could also solve outstanding puzzles of standard big-bang cosmology, such as why the Universe is, to a very good approximation, (=) flat and isotropic (i.e., it looks the same in all directions). |
| (1475) Yet we do not yet have a compelling model, based on (e) fundamental particle physics principles, that explains inflation. |
| (1476) And despite its explanatory power and a great deal of suggestive evidence, we still lack an unambiguous and direct probe of (e) inflation. |
| (1477) Theorists have developed different models for inflation, which all share (e&eb) a common, robust prediction: Inflation would have created (eb) a background of gravitational waves that could have (eb) an observable effect. |
| (1478) These waves would cause (eb) subtle, characteristic distortions of (e) the cosmic microwave background (CMB)—the oldest light in the Universe, released when (e) photons decoupled from matter and the Universe became (=) transparent to radiation |
| (1479) Until now, such a background of gravitational waves has (e) escaped detection. It is therefore very exciting that the BICEP2 (Background Imaging of Cosmic Extragalactic Polarization) collaboration has now reported in Physical Review Letters the observation of polarization modes of CMB radiation [1] that are consistent with (e&eb) such a background. |
| (1480) If confirmed, BICEP2's finding would provide (eb) the first direct probe of physical processes that occurred a mere 10−3610-36 seconds after (e) the big bang. |
| (1481) The result also offers a rare opportunity to directly test theoretical ideas, including inflation and grand unification (the theory that merges the electromagnetic, weak, and strong forces into a single force), at energies 13 orders of magnitude greater than (e) we can currently access at the Large Hadron Collider. |
| (1482) Inflation is based on (e) the possibility that a scalar field in the early Universe got momentarily stuck in (eb) a "false vacuum state"—a metastable condition in which the Universe wasn't at (eb) its global potential minimum. |
| (1483) The energy stored in this vacuum state could drive (eb) an exponential expansion of spacetime (similar to, but much faster than, the expansion due to dark energy seen today) before (e) a phase transition to the true vacuum state occurs, releasing (eb) this energy into (e&eb) particles and radiation. |
| (1484) This creates (eb) an initially hot state for (e) a subsequent big-bang-like expansion from (e) which our Universe subsequently evolved [2]. |
| (1485) If the Universe grew by at least a factor of 10261026, (eb) inflation can explain (eb) both the apparent flatness and isotropy of the Universe [2]: |
| (1486) The expansion would stretch (eb+ eb-dot) any initial curvature of (e) the Universe to near flatness, and smooth out (e&eb) any initial inhomogeneities, just as stretching (eb+ ebdot) a piece of fabric would do. |
| (1487) Isotropy is justified by the fact that the Universe, before inflation, was once much smaller and could have been contained in (eb) a causally connected region that was in thermal equilibrium, resulting in (eb) the approximately uniform, isotropic temperature distribution observed on (eb) large scales today. |

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| (1488) Importantly, the theory can explain the (eb) origin of structure in the Universe: Microscopic quantum fluctuations associated with (e&eb) the scalar field driving inflation could have been "frozen in" during the exponential expansion, resulting, (eb) after inflation, in density fluctuations that might have seeded all observed large-scale structures in the Universe, including (e) all visible galaxies and clusters of galaxies. |
| (1489) The signs of such primordial fluctuations have indeed been detected through (e&eb) their effect on (e&eb) the CMB, first by the Cosmic Background Explorer (COBE) [3] and later by other missions. |
| (1490) These observational studies detected small temperature variations (anisotropies) in (eb) the CMB and found a spectrum compatible with (e&eb) that predicted by inflation. |
| (1491) The magnitude of these fluctuations suggests (eb) an energy scale of inflation (loosely speaking, the energy stored throughout space in (eb) the scalar field potential that drove inflation) close to (e) 10161016 giga-electron-volts (GeV). |
| (1492) But while consistent with inflation, such fluctuations are not (e) conclusive evidence: Their predicted shape and magnitude is known to vary with (e&eb) the details of the inflationary potential (whose form is not yet experimentally constrained). |
| (1493) Hence many different observations could be accommodated, and any agreement between theory and observation is at best suggestive. There is also a possibility that other, noninflationary sources might contribute. However, gravitational waves—the propagating ripples in spacetime predicted by Einstein in 1916—might provide (eb) direct evidence of inflation. |
| (1494) During inflation, gravitational fields would also fluctuate because of (e) quantum effects. |
| (1495) These fluctuations would similarly get "frozen in" and re-emerge after inflation as (=) classical gravitational waves, whose amplitude depends only on (=) the energy stored in the field driving (eb) inflation, and not on other model-dependent details [4]. |
| (1496) If the scale of inflation were indeed as high as a few times 10161016 GeV, this background could produce (eb) specific temperature anisotropies in the CMB [5]. |
| (1497) These anisotropies have thus far not been detected by (e) the most recent searches by the WMAP and PLANCK satellites, and their absence was used to (e) put upper limits on the scale of inflation somewhat below 10161016 GeV. |
| (1498) But a more sensitive CMB probe of (e) gravitational waves exists. |
| (1499) Large-wavelength gravitational waves produced by (e) inflation would have alternatively compressed and expanded (e&eb) space, producing (eb) spatial variations in (eb) the intensity of radiation bombarding free electrons in (eb)the early Universe (see, e.g., Ref: [6]). |
| (1500) As photons decoupled from (e) matter and CMB radiation was released, scattering from (e) the undulating sea of electrons would have created (eb) patterns in the polarization of CMB. |
| (1501) Such polarization anisotropies are predicted in (eb) inflationary models to be (=) most visible when (e) looking at the sky on spatial scales comparable to (e&eb) or somewhat larger than the size of (e) the Universe's horizon at (eb) the time when the CMB was created. |
| (1502) This corresponds to (e&eb) about 1 degree of the field of view of the sky. |
| (1503) Over the past decade, various CMB polarization experiments have thus chased the observation of polarization anisotropies on (e&eb) these scales but, so far, also yielded (eb) merely upper limits on a possible gravitational background. |
| (1504) All of this has now changed with (e&eb) the BICEP2 collaboration's [1] report of (e) a signal that has the predicted (eb) characteristics of such a background at precisely the angular scales where (e) the inflationary signal is expected (see Fig. 1, top deleted due to spatial restraints). |

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| (1505) The BICEP2 Collaboration uses (e) a microwave telescope located at the South Pole, with (e&eb) an array of polarization-sensitive detectors pointing at (eb) a small field (a solid angle of about 11 percent of the full sky) with a resolution of approximately 0.50.5 degrees, thus sufficient to probe the relevant angular scales for the detection of (e) polarization anisotropies originating from (e) **inflation**. |
| (1506) BICEP2's search is based on (e) two expected features of primordial gravitational waves. |
| (1507) First, gravitational waves would produce (eb) a characteristic spatial polarization pattern on CMB—a consequence of (e) their quadrupole nature. |
| (1508) One can decompose (e&eb) any spatial polarization pattern into (e&eb) two orthogonal components, a "curl-free" component and (e&eb) a component involving a twisting pattern. |
| (1509) In analogy to electromagnetic modes, these polarization patterns are called "EE modes" and (e&eb) "BB modes," respectively (see Fig. 1). |
| (1510) Scalar fluctuations in the density of (e) matter would produce (eb) EE modes, while (e) gravitational waves would instead generate (eb) a BB-mode pattern. |
| (1511) The second feature is the aforementioned characteristic spatial scale of (e) expected CMB polarization signatures, including (e) wavelength scales comparable to (e&eb) or larger than the horizon scale at the time the CMB was created. |
| (1512) Since inflation would produce (eb) both gravitational waves and (e&eb) scalar density perturbations, BICEP2, like all CMB-polarization experiments, quote their results in terms of (e&eb) the ratio rr of the amplitude of (e) tensor perturbations (associated with (e&eb) gravitational waves) to (e) the amplitude associated with (e&eb) scalar density fluctuations (reflected in (eb) the magnitude of CMB temperature fluctuations measured by (e) experiments like COBE). |
| (1513) A nonzero rr implies (eb) the detection of a BB-mode pattern and primarily measures (eb) the strength of any possible gravitational wave signal which, if (e) it comes from inflation, is then determined by (e) its energy scale (e.g., see Ref. [7] for a review). |
| (1514) BICEP2 reported a nonzero value of r=0.2(+0.07−0.05)r=0.2(+0.07-0.05), on scales ranging from (e) the horizon scale up to 4–5 times larger, with (e&eb) a claimed significance of better than 5σ5σ. |
| (1515) If confirmed, a value of 0.2 corresponds to (e&eb) an energy density during (e&eb) inflation of ∼1.2×1016∼1.2×1016 GeV. |
| (1516) The claimed signal is far larger than expected, partly because (e) the PLANCK satellite, despite not having yet reported its polarization measurements, has (e) previously put an upper limit on (eb) rr (<0.11<0.11). PLANCK's estimate is based on (e) the absence of a gravitational wave signature on (eb) the CMB temperature anisotropy on (eb) its largest scales (PLANCK observes the whole sky and therefore angular scales up to 360360 degrees). |
| (1517) The apparent discrepancy between the PLANCK upper bound and (e&eb) the BICEP2 detection might be a cause for concern. |
| (1518) However, the two bounds probe different spatial scales, and the BICEP collaboration has noted that, while tensor perturbations are expected to be (=) roughly scale independent, some scale dependence of (e) scalar density perturbations could relax (e&eb) the PLANCK constraint and bring (eb) the two results into agreement. |
| (1519) Of greater concern, of course, is the possibility that the signal could be determined by (e) foreground contamination—any source along (e&eb) the line of sight of the detector that could distort or contaminate (e&eb) the primordial CMB signal. |
| (1520) BICEP2 tested for (e) a wide variety of possible foregrounds, claiming (eb) they cannot be responsible for (e) the observed signal. |

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| (1521) One obvious candidate arises from (e) gravitational lensing: The intervening mass of (e) astrophysical bodies could distort (e&eb) the spatial pattern of polarization in a way that that could transform (e&eb) an EE-mode signal into (e&eb) a BB-mode signal. |
| (1522) The BICEP collaboration has analyzed this possibility extensively. Based on (e) measurements of (e) CMB temperature (and hence density perturbations) and on (e&eb) models of the mass distribution, they argue that such effects would only be (=) significant on smaller angular scales than those contributing to their measurements of rr. |
| (1523) But since BICEP2's first announcement on March 17, 2014, a number of scientists have raised another foreground concern due to (e) the possible impact of (e&eb) our Galaxy's dust. |
| (1524) BICEP2 modeled dust contamination scenarios based on (e) a variety of theoretical models, using data from (e) its predecessor BICEP1 and preliminary dust estimates from (e) the PLANCK satellite. |
| (1525) They concluded that dust could not (e) reproduce the magnitude of the observed signal. |
| (1526) However, PLANCK has in the meanwhile released (eb) more complete measurements of variation in polarized dust foregrounds in (eb) many regions of the Galaxy. |
| (1527) Their analysis increases the uncertainty associated with (e&eb) the dust contribution to (e) the BICEP2 result and also accommodates (e) a potentially larger polarized dust contribution. |
| (1528) Some researchers argued that dust polarization at a level 2–3 times higher than that assumed by BICEP2 is possible and could therefore plausibly be responsible for (e) the entire signal (e.g., see Ref: [8]). |
| (1529) While the BICEP2 collaboration continues to maintain that this is extremely unlikely based on (e) their estimates, many other researchers are more skeptical. |
| (1530) In any case, a variety of upcoming measurements, from PLANCK and other detectors are (=) expected to resolve this important ambiguity within the coming year. |
| (1531) Even if the BICEP2 signal does arise from (e) primordial gravitational waves, various consistency checks will be required (e) before it can be definitively argued that the signal arises from (e) inflation. |
| (1532) Inflation is certainly the best-motivated possibility, and the spectral shape of the signal, at least for (e) the five largest-scale modes observed by BICEP2, agrees extremely well with (=) inflationary predictions (see Fig. 1, bottom). |
| (1533) Several results have already appeared that strongly constrain (e) other possible sources, ruling out (e) for example, the possibility of gravitational waves generated by (e) other possible phase transitions after inflation [9] |
| (1534) . This conclusion is reinforced by more careful analyses of the scale dependence of (e) the predicted signal (e.g., Ref. [10]). |
| (1535) But it will nevertheless be important to assure that all possible alternative sources of (e) gravitational waves are ruled out. |
| (1536) If it can be established that the BICEP2 represents (eb) a gravitational wave signal from (e) inflation, there is one additional exciting implication. |
| (1537) Standard calculations suggest (eb) that the gravitational wave signal from inflation is due to (e) quantum fluctuations in the gravitational field, and this appears to require (e) the existence of gravitons—the hypothetical elementary particles that mediate (e&eb) the force of gravity. |
| (1538) Research by Frank Wilczek and myself, based on (e) a dimensional analysis argument [11], suggests that,(eb) independently of the method used to calculate the gravitational wave spectrum, a quantum gravitational origin is (=) required. |
| (1539) If so, the BICEP2 result then implies (eb) that gravity is ultimately a quantum theory, a result of fundamental importance for physics. |

(1540) It is hard to imagine a more important or wider-ranging cosmological observation. If confirmed by upcoming experiments, the BICEP2 results will push the frontiers of empirical physics forward in an unprecedented way, offering (eb) the first direct observation of gravitational waves and of quantum gravity effects as well as a new probe of the earliest moments in the Universe's life and of physics at grand-unified energy scales. This will thus constrain both our understanding of fundamental interactions and the evolution of our Universe as a whole.

**Detection of B-Mode Polarization at Degree Angular Scales by BICEP2 P. A. R. Ade et al. (BICEP2 Collaboration) Phys. Rev. Lett 112, 241101 – Published 19 June 2014**

(1541) Authors report results from the BICEP2 experiment, a cosmic microwave background (CMB) polarimeter specifically designed to search for (e) the signal of inflationary gravitational waves in (eb) the B-mode power spectrum around $\ell \sim 80$.

(1542) The telescope comprised a 26 cm aperture all-cold refracting optical system equipped with (e&eb) a focal plane of 512 antenna coupled transition edge sensor 150 GHz bolometers each with temperature sensitivity of $\approx 300\ \mu K_{CMB}\sqrt{s}$. BICEP2 observed from (e) the South Pole for three seasons from 2010 to 2012.

(1543) A low-foreground region of sky with (e&eb) an effective area of 380 square deg was observed to (e) a depth of 87 nK deg in Stokes Q and U.

(1544) In this paper authors describe the observations, data reduction, maps, simulations, and results. They find an excess of B-mode power over the base lensed-ΛCDM expectation in the range $30<\ell<150$, inconsistent with (e&eb) the null hypothesis at (eb) a significance of $>5\sigma$.

(1545) Through jackknife tests and simulations based on (e) detailed calibration measurements authors show that systematic contamination is (=) much smaller than the observed excess.

(1546) Cross correlating against WMAP 23 GHz maps we find that Galactic synchrotron makes a negligible contribution to the observed signal. We also examine a number of available models of polarized dust emission and find that at their default parameter values they predict power ~(5–10)× smaller than the observed excess signal (with no significant cross-correlation with our maps). However, these models are not sufficiently constrained by external public data to exclude the possibility of dust emission bright enough to explain the entire excess signal. Cross correlating BICEP2 against 100 GHz maps from the BICEP1 experiment, the excess signal is confirmed with $3\sigma$ significance and its spectral index is found to be consistent with that of the CMB, disfavoring dust at $1.7\sigma$. The observed B-mode power spectrum is well fit by a lensed-ΛCDM+tensor theoretical model with tensor-to-scalar ratio $r=0.20+0.07-0.05$, with $r=0$ disfavored at $7.0\sigma$. Accounting for the contribution of foreground, dust will shift this value downward by an amount which will be better constrained with upcoming data sets. Received 4 April 2014 DOI:http://dx.doi.org/10.1103/PhysRevLett.112.241101 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**(1547) Physicist claims to have observed quantum effects of Hawking radiation in the lab for the first time August 16, 2016 by Bob Yirka report**

(1548) This computer-simulated image shows a supermassive black hole at the core of a galaxy. The black region in the center represents the black hole's event horizon, where no light can escape the massive object's gravitational grip. The black …more

(1549) (Phys.org)—Jeff Steinhauer, a physicist at the Israel Institute of Technology, has published a paper in the journal Nature Physics describing experiments in which he attempted to create a virtual black hole in the lab in order to prove that Stephen Hawking's theory of radiation emanating from black holes is correct —though his experiments are based on sound, rather than light. In his paper, he claims to have observed the quantum effects of Hawking radiation in his lab as part of a virtual black hole—which, if proven to be true, will be the first time it has ever been achieved.

(1550) For many years, scientists believed that nothing could ever escape from a black hole. But in 1974, Stephen Hawking published a paper suggesting that something could—particles that are now called Hawking radiation. His idea was that if a particle (and its antimatter mate) appeared spontaneously at the edge of a black hole, one of the pair might be pulled into the black hole while the other escaped, taking some of the energy from the black hole with it—which would explain why black holes grow smaller and eventually disappear. Because such emissions are so feeble, no one has been able to measure Hawking radiation, so researchers have instead tried to build virtual black holes in labs to test the theory. One type of virtual black hole was proposed back in 1981 by Bill Unruh with the University of British Columbia—he suggested that an analogue might be created using water instead of light. He imagined a phonon existing at the edge of a waterfall—as the water speeds up, it begins to move faster than the speed of sound, causing it to be trapped. But if the phonon had an entangled mate that eluded the fall by moving away before getting caught up, it could escape. In this new effort, Steinhauer has built a device based on that idea and in so doing, claims he has observed an analogue of Hawking radiation.

(1551) The experiment consisted of creating an entangled pair of phonons sitting inside a bit of liquid that had been forced (via laser) to move very fast and then observing the action as one of the pair was pulled away as the liquid began to move faster than the speed of sound, while the other escaped—the fluid was a Bose-Einstein condensate of rubidium-87 atoms. After repeating the experiment 4,600 times Steinhauer became convinced that the particles were entangled, a necessity for a Hawking radiation analogue. His findings do not prove Hawking's theory to be true, of course, but they do appear to add a degree of credence that other researchers have thus far not been able to achieve. More information: Jeff Steinhauer. Observation of quantum Hawking radiation and its entanglement in an analogue black hole, Nature Physics (2016). DOI: 10.1038/nphys3863

**(1552) Abstract**

We observe spontaneous Hawking radiation, stimulated by quantum vacuum fluctuations, emanating from an analogue black hole in an atomic Bose–Einstein condensate. Correlations are observed between the Hawking particles outside the black hole and the partner particles inside. These correlations indicate an approximately thermal distribution of Hawking radiation. We find that the high-energy pairs are entangled, while the low-energy pairs are not, within the reasonable assumption that excitations with different frequencies are not correlated. The entanglement verifies the quantum nature of the Hawking radiation. The results are consistent with a driven oscillation experiment and a numerical simulation. Journal reference: Nature Physics search and more info website

Read more at: http://phys.org/news/2016-08-physicist-quantum-effects-hawking-lab.html#jCp

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as

the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t - T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

These equations can be written for

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(10) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.

(11) However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].

(12) Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.

Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

**TIME DELAY**

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period,

which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types system takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t - T_7)F_7(R_7(t - T_7) - N_7(t)M_7((R_7(t) \qquad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7}+\frac{N_7 dG_7}{dR_7}-\frac{dB_7}{dR_7}>\frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7}+\frac{dM_7}{dR_7}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also normalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_7$) and the other derived from the non-delayed or "immediate" processes($I_7$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_7=\frac{N_7 G_7}{2M_7}\frac{d(7nF_7)}{dR_7}(\geq 0) \quad (7a)$$

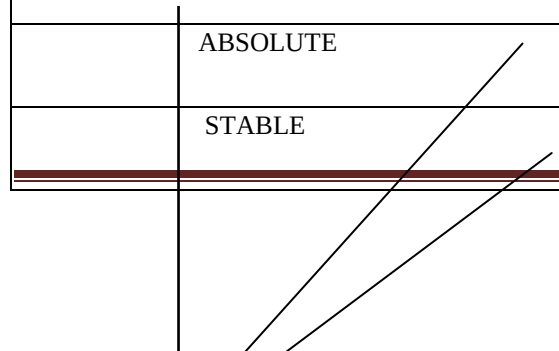
$$I_7=-\frac{N_7 G_7}{2M_7}\frac{d(7nM_7}{dR_7}(\geq 0) \quad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7>L_7-I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

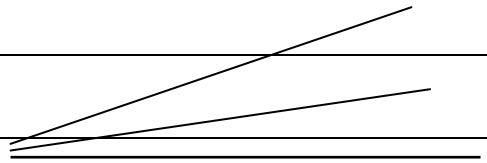
$$P_7>L_7-I_7, \quad \text{when} \quad P_7^2/4>I \quad (8a)$$

$$P_7>L_7-I_7+\frac{(I_7-P_7^2}{(L_7-I)} \quad \text{when} \quad P_7^2/4<I. \quad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,





**DELAYED PREDATOR CREDIT TYPES EFFECT WITH ITS WIDE RANGING RAMIFICATION**

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

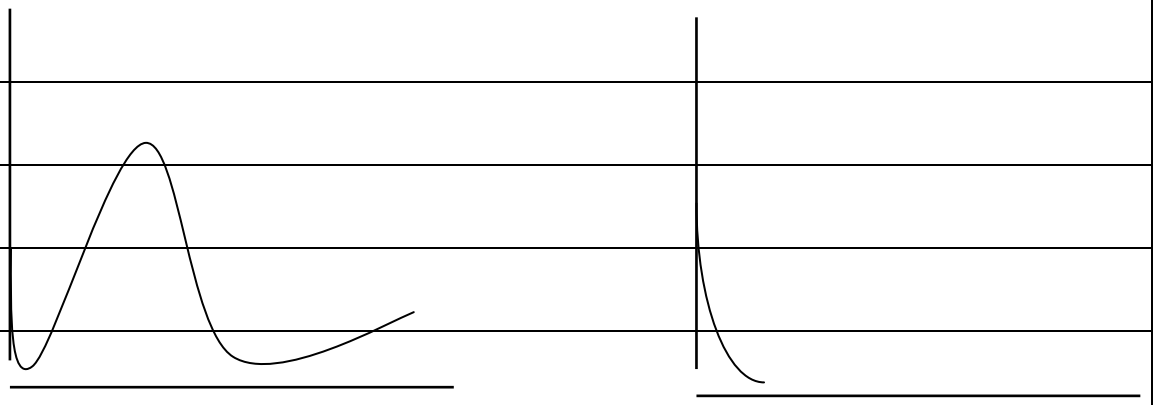
It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay

and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



**TIME DELAY**

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

To scores of problems this relationship is important. Of them some of the examples are as follows:

**New look at the galaxy-shaping power of black holes**

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| (385) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving. |
| (386) "These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science. |
| (387) "Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy." |
| (388) This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return. |
| (389) The remaining matter gets absorbed into the black hole, adding to (e) its mass. |
| (390) The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne. |
| (391) The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution. |
| (392) Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form. |
| (393) "Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara. |
| (394) "It creates (eb) these bubbles, and the bubbles keep the plasma hot. |
| (395) That's what prevents galaxies from (e) becoming even bigger than they are now." |
| (396) Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space. |
| (397) In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene. |
| (398) Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve. |
| (399) "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara. |
| (400) "It's like nature's thermostat, which keeps these galaxies from (e) growing. |
| (401) If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars." |

**A new look at the galaxy-shaping power of black holes**

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| (402) An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the |
| (403) McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said. |

(404) Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.

**Doomed satellite missions**

(405) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.

(406) At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.

(407) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.

(408) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.

(409) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.

(410) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people

(411) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."

(412) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.

(413) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(414) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.

(415) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.

(416) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.

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| (417) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses. |
| (418) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well. |
| (419) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma. |
| (420) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6. |
| (421) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus. |
| (422) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (423) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (424) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (425) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (426) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (427) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (428) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (429) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (430) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (431) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (432) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (433) We know dark matter exists only from (e) its gravitational effects. |
| (434) It holds together galaxies and (e&eb) galaxy clusters. |

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| (435) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (436) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |
| (437) They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (438) X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (439) They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (440) This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (441) An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |
| (442) Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (443) "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (444) The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (445) Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (446) Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (447) "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (448) Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (449) "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New |

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| York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (450) Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories. |
| (451) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen. |
| (452) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer. |
| (453) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants. |
| (454) These were treated as (=) "aware" and "unaware" trials, respectively. |
| (455) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it. |
| (456) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another. |
| (457) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion. |
| (458) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness. |
| (459) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said. |
| (460) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness." |
| (461) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner. |
| (462) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain. |
| (463) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said. |
| (464) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

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| (465) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter. |
| (466) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined. |
| (467) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions. |
| (468) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda. |
| (469) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general. |
| (470) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole. |
| (471) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes. |
| (472) Noda and his colleagues instead studied the change in emissions over (e&eb) time. |
| (473) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection. |
| (474) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1). |
| (475) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole. |
| (476) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved. |
| (477) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp |

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

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| (478) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking “palm tree” organization. |
| (479) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society |
| (480) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible? |
| (481) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem. |
| (482) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other. |
| (483) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible. |
| (484) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges. |
| (485) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century. |
| (486) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching. |
| (487) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom. |
| (488) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth. |

**Spheres on trees**

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| (489) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration. |
| (490) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement. |
| (491) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum. |
| (492) By visualizing the problem in this way, the researchers noticed that these particular graphs don’t (e) have lots of branches extending from lots of (e) other branches. |

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| (493) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure. |
| (494) The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network. |
| (495) As a result, it provides (eb) an important clue for finding the global minimum. |
| (496) It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum. |
| (497) This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large. |

**Thompson's problem 2**

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| (498) By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society |
| (499) The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima. |
| **(500)** The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node. |

**Implications for Thomson's and Smale's problems**

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| (501) Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before. |
| (502) "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around." |
| (503) As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem. |
| (504) "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view." |
| (505) In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges. |
| (506) From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear. |
| (507) "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK. |
| (508) "For larger numbers, there will likely be (=) multiple funnels. |
| (509) Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum." |

(510) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.

(511) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.

(512) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t-T_7)F_7\big(R_7(t-T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t-T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) and the effective prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this

rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Look at the resemblance of String Theory and Abhinava's earlier theory and how the model could be applied: The world of manifestation and differentiated perceptions (Yikalpa) thus extends from one Centre to (eb) the next. **What is meant again is the realisation or the evolution of individual consciousness or general ledger to be (=)the cosmic general ledger is the sole criterion for the Satchidananda state where all dualities cease** Although it is never in fact divorced from (e) the subject who resides there, the ignorant fail (e) to grasp this fact and so, cut off from (e) the Centre, the world of objectivity becomes (=) for them the sphere of Maya. Bhagavatotpala quotes the Light of Consciousness amvitprakdsa): This ever pure experience (suddhanubhava) is variegated by (e) each form [revealed within it] Even so it remains unstained (nirmala) when moving to another. Just as a cloth which is naturally white, once dyed, cannot (e) change colour without [first] becoming white again, similarly the pure power of awareness, (citi) once coloured by (e) form, is pure [again] at the Centre where that form is (=) abandoned and from whence it proceeds to (e&eb) another. In his Essence of Vibration (Spandasarpdoha), K§emaraja explains that the rise and fall of every individual perception in the field of awareness is (=) a specific pulsation of consciousness. This is exactly the individual general ledger transaction lifted to the transaction in cosmic general ledger. During The Initial Instant Of Perception, T Consciousness Is Manifestly Apparent And The Yogi, Participating In Its Plenitude, Observes The Outer World Without Being Attached To Any Particular Divine Body And Sacred Circle Of The Senses Singling It Out From Any Other, Like A Man Who Observes A City From A High Mountain Peak.

He Sees The Outer World Reflected Within His Consciousness Free Of Thought-Constructs And So 'Stamps* The Outer On The Inner While Absorbing The Object And Means Of Knowledge In The Pure Subject Which Grasps Them As The Expansion Of His Own Nature. K§emaraja Says: By Penetrating Into Bhairavimudra, The Yogi Observes The Vast Totality Of Beings Rising From, And Dissolving Into, The Sky [Of Conscious- Ness], Like A Series Of Reflections Appearing And Disappearing Inside A Mirror. Through (e&eb) The Practice Of Bhairavimudra The Yogi Realises That He Is (=) The Substratum Consciousness (Adhisfhatr) Which Both Underlies (e&eb) And Is (=) The Essence Of All Things. He Discovers That Phenomena Have (e) No Independent Existence Apart From (e) Him And So Are, In This Sense, Void. At The Same Time, He Realises That Because (e) All Things Are Consciousness, They Are (=) Far From Unreal. He Views The Outer World Yet Sees (e) It Not. Beyond Both Voidness And Non-Voidness He Penetrates Into (e&eb) The Supreme Abode (Pararn Padam) Of Siva's Consciousness. [The Powers Of The Senses] Endowed With (e&eb) The Attributes Of The Great Union [Between Subject And (e&eb) Object] Whose Form Is The Wakening Of Man's Spiritual Potential (Kuntfalinf), Fill (e&eb) [With Consciousness] The Outer Clatter Of Diversity (Bhedatfambara) Born Of (e) Its Intense Power And Are Then Established In (eb) The Unobscured Abode Of The Void Of Consciousness To Shine [There] Eternally. Thus Residing Beyond (e) Being And Non-Being, The Sole Protector Of The Unity Which Is Tranquil And Expanding [Consciousness], Whose Glory Is All-Embracing And Form Unobscured, Is Called Bhairavimudra. Through (e&eb) The Practice Of Bhairavimudra, The Yogi Unites The Universal Vibration Of T Consciousness With (e&eb) The Individual Pulsation Of Objectivised 'This' Consciousness. The Two Aspects Of Consciousness Are Now In (eb) A State Of Equilibrium Like (=) The Two Pans Of An Evenly Weighed Balance And The Yogi Experiences (e&eb) The Pure Knowledge (Suddhavidya) That: 'I Alone Am (=)All Things'. Thus Becoming (=) The Master Of The Wheel Of Energies He Is (=) Free, Like Siva, To Create And Destroy. When [The Yogi] Is Well Established, Without (e) Wavering, Solely In The Integral Egoity Of His Authentic Nature, The Spanda Principle, And Is

Absorbed In Contemplation (Samavisfa), He Becomes (=) One With It The Doctrine Of Vibration (Tanmaya). Then . . . Dissolving And Creating The Universe By Means Of (er) His Introverted And Extroverted Absorption, He Destroys And Creates (e&eb) All Things Out Of Sarikara, His Innate Nature. [Thus] He Assumes The State Of The Universal Experiencer And Having Absorbed All That Is To Be Experienced From [The Grossest Level] — Earth — To (The Subtlest) — Siva — He Reaches (eb) The State Of The Supreme Subject By (e) Progressively Recognising [His Identity With Him]. Thus, Introverted And Extroverted Absorption Both Lead To (eb) The Recognition Of The Pulsation (Spanda) Of (e) One's Own Consciousness. This Is Where (e) Individual General Ledger Becomes (=) Cosmic General Ledger. At The Level Of Consciousness Corresponding To (e&eb) Siva's Basic State (Saryibhavdvastha), The Alternation From Inner To (e&eb) Outer Is Instantaneously Resolved Into (e&eb) The Vibration Of His Nature. When The Yogi Finally Comes To Be (=) Constantly Aware Of This Reality, His Enlightenment Is (=) Full And Perfect. Freed Of All Means (Anupaya) And Delighting In The Power Of His Bliss (Anandasakti), He Knows And Does Whatever He Pleases. The Yogi Seeking Self-Realisation Must Acquire (e) Mastery Over This Movement. Ksemaraja Stresses That The Doctrine Of Vibration Teaches That Liberation Can Only Be Achieved By (e) First Withdrawing All Sense Activity In Introverted Contemplation (Nimilanasamadhi) To (e) Then Experience The 'Great Expansion' (Mahdvikdsd) Of (e) Consciousness While (e) Recognising This To Be (=) A Spontaneous Process Within (eb) It. This Is Done Through (e&eb) The Practice Of Kramamudra. A Passage From The Now Lost Kramasiitra Explains: Although The Adept's Attention [May Be] Outwardly Directed, (eb) He Enjoys Contemplative Absorption Through (e&eb) The Introverted Aspect Of (e) Kramamudra. Initially He Turns Inward From (e)The Outside World And [Then] From Within [Himself] He Exits Into The Outer World Under The Influence Of His Absorption. There Is No Necessity Of Any Glorification Or Mortification Here. It Is Just Like Joining A Job And Retiring And Realising All That Happened Is Just By His Own Effort And Nothing Else. Thus The Sequence (Krama) In This Attitude (Mudra) [Ranges Through] Both Inner And Outer. The Yogi Must Pervade

The Surface Level Of Awareness (Yyutthana) With The Same Bliss He Experiences Plunging Into The Depths Of Con- Templative Absorption (Samadhi). Note Here That Samadhi Is A Death Like State And Still Being Matter And Information Stored You Experience Consciousness. Submerging Himself And Emerging Repeatedly From Samadhi, He Eventually Recognises That The Unity Of Consciousness Pervades Both States: The Best Of Yogi's, Who Has Achieved A State Of Complete Absorption Even When Risen From Meditation, [Inwardly] Vibrating Like A Drunkard In Blissful Inebriation From The After-Effects Of The Nectar Of Con- Templation, Sees All Things Dissolving In The Sky Of Consciousness Like A Cloud In The Autumn Sky. He Plunges Repeatedly within Himself Divine Body and Sacred Circle of the Senses and Becomes (=) Aware Of His Identity With (e&eb) Consciousness By (e) The Practice Of Introverted Contemplation. Thus Even When He Is Said To Have Risen From (e) Absorption, He Is (=) One With [His] Experience Of It. Ksemaraja Goes On To Explain That This Practice Is Called 'Mudra' Because (e) It Both Fills The Adept With (e&eb) Bliss (Mud) And Is Itself The Bliss Of (e) Consciousness. Moreover, It Dissolves (e) Away (Dra) All Bondage And 'Stamps' The Universe Of Experience With (e&eb)The Seal (Mudra) Of The Fourth State (Turiya) Of Enlightened Consciousness Beyond, And Including, The Three States Of Waking, Dreaming And Deep Sleep. It Is Called 'Krama' Because (e) It Is The Root Source Of (e) All Emanation And All Other Conscious Processes Which Succeed (e&eb) One Another In Ordered Sequence (Krama) And Is, (=)At The Same Time, Their Successive (Krama) Appearance As Well. By The Practice Of Kramamudra The Opposites Fuse And Siva And Sakti Unite. They Yogi Comes To Experience The Simultaneous Pervasion Of All The Lower, Grosser Categories Of Existence By The Higher And The Presence Of The Lower In The Higher. Commencing His Practice In A Low Form Of Bhairavimudra, The Yogi Conjoins The Outer With The Inner; Then, In Kramamudra, He Fills Both The Outer With (e&eb) The Inner And (e&eb) The Inner With (e&eb)The Outer. When He Achieves (eb) Perfection In This Two-Fold Movement, He Attains To (e) The Highest Form Of Bhairavimudra In Which The Two Merge Completely In (eb)The Experience Of (e) The Absolute (Anuttara), Free

Of (e) All Differentiation And (e&eb) Polarities. If He Fails To Maintain Awareness Of This State, He Again Falls Into (eb) Kramamudra Until He Has Finally Completely Merged All The Highest States In (e&eb) The Lower And The Lower In (e&eb) The Higher. He Then No Longer Needs To (e) Resort To (e) Any Means (Anupaya) To (e) Achieve Liberation. All He Says or Does Anything He Perceives or Thinks, Instantly Occasions In (eb) Him The Highest Level Of Consciousness. Thus The Fruit Of Bhairavimudra Is (*=)The Wonder (Camatkara) Or Amazement (Vismaya) That Overcomes (e)The Yogi When (e) He Reaches The Plane Of (e) Union (Yogabhumika) Y Xsl Where (e) All Opposites Merge In (e&eb)The Radiance Of The Great Light Of Consciousness. The Stanzas On Vibration Teach: How Can One Who, As If Astonished, Beholds (e) His Own Nature As (=) That Which Sustains [The Existence Of Everything] Be Subject To (e&eb) This Painful Round Of Transmigration (Kusrti)? The Yogi, Recognising His True Nature To Be (=)The Supreme Subject, Is Astonished To (e) Suddenly Discover That The Individual He Thought He Was, Caught Up (e&eb) In The Trammels Of Thought And Living In A World Enmeshed In (e&eb) The Web Of Time And Space, Does Not Really Exist At All. He Experiences A 'Turning About' (Paravrtti) In (eb)The Deepest Seat Of The Doctrine Of Vibration Consciousness As He Penetrates His True Nature. The Sudden Eruption Of This Intuition (Pratibha) Arouses In Him A Cry Of Amazement As He Transcends All Thought-Constructs And, Perfectly Absorbed In His Own Nature, Is Liberated. The Path To Liberation Essentially, Spanda Doctrine Is Concerned With Two Matters. The First Is To Impart To Those Who Are Fit To Receive The Teachings A Deeper Understanding Of The Ultimate Goal Of Life (Upeya). When We Have Understood What Truly Benefits Us And Is Worth Attaining And What, On The Contrary, Is Of No Real Value But Stands In The Way Of This Attainment, We Can Begin To Make Progress Towards Our Goal. This Is Spanda Doctrine's Second Concern, Namely, To Show The Way In Which We Can Develop Spiritually Through Siva's Grace And The Right Application Of The Means To Realisation That It Teaches. When Both These Aspects Of The Teaching Have Been Correctly Understood And Applied, The Spanda Yogi Achieves A Clear And Permanent

Realisation Of His Goal And Is Liberated, Thus Fulfilling The Ultimate Aim Of The Teaching. The Doctrine Of Vibration Is Not Meant For The Spiritually Dull. It Is Not For The Worldly Whose Consciousness, Clouded By Ignorance, Is As If Dreaming, Even During The Waking State Of Daily Life, The Dream Of Its Own Thought-Constructs. The Teachings Are Meant For Those Who Are Awake (Prabuddha), Those Who, Full Of Faith And Reverence, Are Always Alert And Intent On Discerning The True Nature Of Ultimate Reality. This Reality Is Understood In Three Basic Ways. The First Is Purely Transcendental. The Stanzas Choose This Aspect As The One Which Formally Defines It Most Specifically. Ultimate Reality Transcends All The Opposites, Including Subject And Object. This Does Not Mean, However, That It Is An Unconscious Void, A Mere Absence Of All Existence. In Fact, This Negative Characterization Of Reality (Which Includes Also A Denial Of All That Is Unconscious) Implies A Positive Immanence In Which The Opposites Are United In The Oneness Of Pure Consciousness That Is Equally Siva And Spanda, His Universal Activity. These Two Seemingly Contrasting Aspects Are Reconciled In The Third, Namely, Reality Understood As The Essential Nature Of All Things. Although Universal And Everywhere The Same, It Is Understood To Be The Essential And Specific Nature Of Each Existent As Its 'Own Nature' (Svabhava).

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such

a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

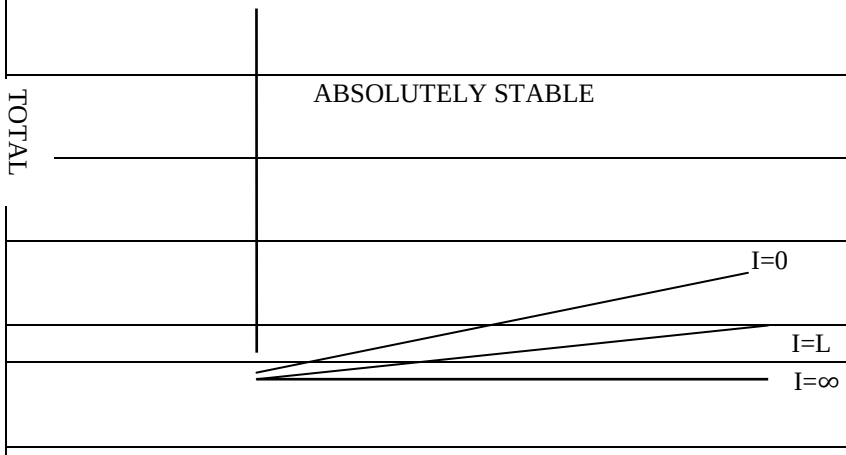
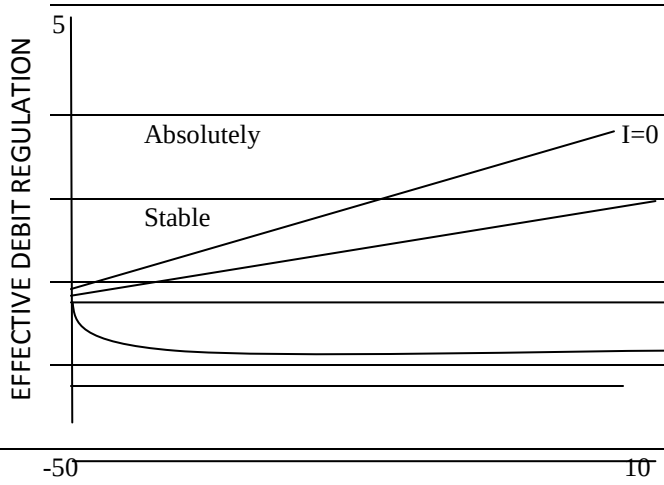



FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

This equation is of paramount importance in the following cases: According to the principles of geometric quantum mechanics, the physical characteristics of a given quantum system can be represented by geometrical features that are preferentially identified in this complex manifold. Here we construct a number of examples of such features as they arise in the state spaces for spin, spin 1, spin and spin 2 systems, and for pairs of spin systems. A study is then undertaken on the geometry of entangled states. A locally invariant measure is assigned to the degree of entanglement of a given state for a general multi-particle system, and the properties of this measure are analysed for the entangled states of a pair of spin particles. With the specification of a quantum Hamiltonian, the resulting Schrödinger trajectories induce an isometry of the Fubini–Study manifold, and hence also an isometry of each of the energy surfaces generated by level values of the expectation of the Hamiltonian. For a generic quantum evolution,

the corresponding Killing trajectory is quasiergodic on a toroidal subspace of the energy surface through the initial state. When a dynamical trajectory is lifted orthogonally to Hilbert space, it induces a geometric phase shift on the wave function. The uncertainty of an observable in a given state is the length of the gradient vector of the level surface of the expectation of the observable in that state; a fact that allows us to calculate higher orders corrections to the Heisenberg relations. A general mixed state is determined by a probability density function on the state space, for which the associated first moment is the density matrix. The advantage of a general state is in its applicability in various attempts to go beyond the standard quantum theory, some of which admit a natural phase-space characterization. **Journal of Geometry and Physics Volume 38, Issue 1, April 2001, Pages 19–53 Geometric quantum mechanics Dorje C. Brody Lane P. Hughston DOI: 10.1016/S0393-0440(00)00052-8** Expansion (unmesa) of [consciousness] or the creative intuition {pratibha) is (=) [experienced] in the interval which divides two [moments] of differentiated perception (vikalpa) It is here that they arise and disappear. The Sastras and Agamas proclaim with reasoned argument that it is (=) free of thought-constructs {Nirvikalpa) and precedes (e) Path to Liberation all mental representations of any object. None can deny that a gap exists between perceptions insofar as two moments of thought are (=) invariably divided. This [gap] is (=) the undifferentiated unity of all the countless manifestations. Similarly, in the outer more objective sphere, where change consists (e) of the alterations in the configurations of manifest appearances (Abhasa\ the transition from one to (e&eb) another corresponds to (e&eb) a phase of pure luminosity that marks (eb) the beginning of one form and the end of another

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^{\infty} \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7 \left( U_7 - V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7 \right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$then given the constraints of $f(T_7), (A1.5) can\ never\ be\ satisfied\ Co$nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7 \lambda_7 e^{-T_7 \lambda_7} + [(U_7 - V_7) M_7 + Y_7] - [((U_7 -$$

V 7M 7–]**e**–7 7λ 7=0

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$\boldsymbol{Z}_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7) \cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7) \sin(Z_7 T_7) = 0$$

The warrant of Truth is ever elusive when we deconstruct the foundational justifications of our convenient notions about the way the world works. It is easy to confuse what is actually the creation of beliefs with the "discovery of Truth," a common goal of science and theology. Science offers no ontological Scientific Picture of the World but substitutes a number of applied theories as meta-theory. It's close, powerful and useful: it 'works' and allows use to do work. But the still mysterious "true" metatheory may reveal the basis of consciousness. But who knows when we will lift that veil? Permeating the living reality of our culture are certain contagious notions, fads and trends that have the ability to influence the way we think about the nature of Reality and ourselves. Some of them are toxic and can consume you; others are just whims of pop culture infectious and benign as the common cold. **Jung** described the concept of psychic contagion by certain archetypal forces inherent in the human psyche, which manifest in our spiritual lives and belief systems. He spoke of both conscious and unconscious contamination. Notions like this range from simple superstitions to scientific concepts, to urban myths. Notions sweep through our culture and insinuate themselves within its fabric, as fads, whether they are "real" or not, they can be influential. They mood alter us, make us feel we belong. An analysis of these notions is useful in distinguishing a common human phenomenon from any potential "alien influence" which may or may not be exerted on us from an unknown source. **Strange Attractors: Transference, Holography, and an Archetype Burke, J. (2003).**

**Strange Attractors: Transference, Holography, and an Archetype (Doctoral dissertation, Pacifica Graduate Institute, 2003).** In consideration to the Nobel Proze for Chemistry this year(2014), a stability nalysis of the proble in to single molexues would be a contemporaneous question that has to be taken cognisace of and given importance to . Determination of molecular structure by geometry optimization became routine only after efficient methods for calculating the first derivatives of the energy with respect to all atomic coordinates became available. Evaluation of the related second derivatives allows the prediction of vibrational frequencies if harmonic motion is estimated. More importantly, it allows for the characterization of stationary points. The frequencies are related to the eigenvalues of the Hessian matrix, which contains second derivatives. If the eigenvalues are all positive, then the frequencies are all real and the stationary point is a local minimum. If one eigenvalue is negative (i.e., an imaginary frequency), then the stationary point is a transition structure. If more than one eigenvalue is negative, then the stationary point is a more complex one, and is usually of little interest. When one of these is found, it is necessary to move the search away from it if the experimenter is looking solely for local minima and transition structures. The total energy is determined by approximate solutions of the time-dependent Schrödinger equation, usually with no relativistic terms included, and by making use of the Born–Oppenheimer approximation, which allows for the separation of electronic and nuclear motions, thereby simplifying the Schrödinger equation. This leads to the evaluation of the total energy as a sum of the electronic energy at fixed nuclei positions and the repulsion energy of the nuclei. Notable exceptions are certain approaches called direct quantum chemistry, which treat electrons and nuclei on a common footing. Density functional methods and semi-empirical methods are variants on the major theme. **Orlando Alvarez** discusses Polyakov's quantization of the string in the presence of a boundary allowing for an arbitrary topology for the world sheet. In addition to the dynamical conformal factor discovered by Polyakov, there are a finite

number of new degrees of freedom if the surface is more complicated than a sphere or a disc. The quantization of the Liouville theory in an arbitrary topology is also discussed. A one-loop calculation shows that the model is renormalizable if one performs a mass renormalization and an additive field renormalization. The renormalization group equations have a perturbative infrared unstable fixed point in all topologies. Copyright © 1983 Published by Elsevier B.V. **Nuclear Physics B Volume 216, Issue 1, 25 April 1983, Pages 125–184 Theory of strings with boundaries: Fluctuations, topology and quantum geometry Orlando Alvarez DOI: 10.1016/0550-3213(83)90490-X** Quantum geometry (the modern loop quantum gravity involving graphs and spin-networks instead of the loops) provides microscopic degrees of freedom that account for black-hole entropy. However, the procedure for state counting used in the literature contains an error and the number of the relevant horizon states is underestimated. In our paper a correct method of counting is presented by **Marcin Domagala and Jerzy Lewandowski.** Results lead to a revision of the literature of the subject. It turns out that the contribution of spins greater than 1/2 to the entropy is not negligible. Hence, the value of the Barbero–Immirzi parameter involved in the spectra of all the geometric and physical operators in this theory is different than previously derived. Also, the conjectured relation between quantum geometry and the black-hole quasi-normal modes should be understood again. **Marcin Domagala and Jerzy Lewandowski 2004 Class. Quantum Grav 21 5233 doi:10.1088/0264-9381/21/22/014 Black-hole entropy from quantum geometry** Quantum geometry predicts that a universe evolves through an inflationary phase at small volume before exiting gracefully into a standard Friedmann phase. This does not require the introduction of additional matter fields with ad hoc potentials; rather, it occurs because of a quantum gravity modification of the kinetic part of ordinary matter Hamiltonians. Authors draw the cognizance of application of the same mechanism that can explain why the present day cosmological acceleration is so tiny. DOI: http://dx.doi.org/10.1103/PhysRevLett.89.261301

**Inflation from Quantum Geometry Phys. Rev. Lett 89, 261301 – Published 12 December 2002 Martin Bojowald** The loop quantum cosmology of the closed isotropic model is studied with special emphasis on a comparison with traditional results obtained in the Wheeler-DeWitt approach by **Martin Bojowald** .Investigation includes the relation of the dynamical initial conditions to boundary conditions such as the no-boundary or the tunneling proposal and a discussion of inflation from quantum cosmology. **DOI: http://dx.doi.org/10.1103/PhysRevD.67.124023 Loop quantum cosmology, boundary proposals, and inflation Phys. Rev. D 67, 124023 – Published 23 June 2003 Martin Bojowald and Kevin Vandersloot** In loop quantum cosmology, the universe avoids a big bang singularity and undergoes an early and short super-inflation phase. During super-inflation, non-perturbative quantum corrections to the dynamics drive an inflaton field up its potential hill, thus setting the initial conditions for standard inflation. **Shinji Tsujikawa et al** show that this effect can raise the inflaton high enough to achieve sufficient e-foldings in the standard inflation era. They also analyse the cosmological perturbations generated when slow-roll is **violated after super-inflation** and show that loop quantum effects can in principle leave an **indirect signature** on the largest scales in the CMB, with some loss of power and running of the spectral index. **Shinji Tsujikawa et al 2004 Class. Quantum Grav 21 5767 doi:10.1088/0264-9381/21/24/006 Loop quantum gravity effects on inflation and the CMB Quantum geometry from quantum information** is the study of **Florian Girelli and Etera R Livine .**Loop quantum gravity defines the quantum states of space geometry as spin networks and describes their evolution in time. We reformulate spin networks in terms of harmonic oscillators and show how the holographic degrees of freedom of the theory are described as matrix models. This allows us to make a link with non-commutative geometry and to look at the issue of the semi-classical limit of loop quantum gravity from a new perspective. This work is thought of as part of a bigger project of describing quantum geometry in quantum information terms.

**Florian Girelli and Etera R Livine 2005 Class. Quantum Grav 22 3295 doi:10.1088/0264-9381/22/16/011 Reconstructing quantum geometry from quantum information: spin networks as harmonic oscillators.** Quantum geometric mechanics is the topic of study by **Dorje C. Brody Lane P. Hughston** establishing a clear relation between the quantum mechanics and quantum geometry and also the entangled states. The manifold of pure quantum states can be regarded as a complex projective space endowed with the unitary-invariant Fubini–Study metric. Note how differential gheometry concepts with Nunney's model could be entangle so that it is feasibly and viably applied to various type of model including consciousness (See Doctrine of Vibration)

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $\boldsymbol{Z}_7^2$ and $Z_7$in (A2.2):

$W_7^2 (X_7^2$-3 $M_7\ A_7\ X_7$- $M_7^2 (X_7$- $Y_7) + 2\ M_7^3\ A_7$+2 $A_7^2\ M_7^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_7$=1-cos ($Z_7 T_7$).

In the absence of a delay the system is stable for $A_7 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7 > (X_7 - Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$A_7 > [\ X_7 - Y_7/2\ M_7$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_7, A_7) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_7 \le 2$ *is* of interest First consider the extreme values

$$Q_7(0, A_7) = A_7^2 (X_7 - Y_7)$$

$$Q_7(2, A_7) = (A_7 + 2M_7)^2 (2A_7 M_7 - (X_7 - Y_7))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7 M_7^2 - 3M_7 X_7\right) + W_7\left(3M_7 A_7^2 - 2A_7 X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7 + Y_7\right)$$

$$=2M_7^2 W_7^2 (M_7 + A_7) + M_7 A_7^2 W_7 + (2A_7 M_7 - X_7 + Y_7)(A_7 W_7 + M_7 W_7^2$$

$$+ (2 - W_7)((2M_7 X_7 + M_7 Y_7) W_7 + A_7 (X_7 + Y_7)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $W_7$touches the $W$–axis only at its minimum value in which case (A2.9) shows that

further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$has no minimum in the range 0≤ $W_7$ ≤2.Defining

$Q_7(W_7, A_7) = a_7 W_7^2 + c_7$ *where* $a_7\, b_7$ *and* $c_7$ Are functions of $A_7$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7 \geq b_7$$

And

$$b_7^2 \geq 4a_7 c_7$$

$$Case(i): Y_7 = 0$$

*Let* $Y_7 = 0$ And try the substitution A= $X_7 = /2M_7$ *in* (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_7 X_7 /$ $2M_7$) is positive for *W<2. Since* $Q_7$ increases with $A_7$(see A2.9), increasing $A_7$will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

$Case\ (ii): Y_7 = y_7\ X_7$ ,

$Let\ Y_7 = y_7\ X_7$And try the substitution $A_7 = X_7/(2 + A_7\ y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + n_7 y_7)$ we have

$$W_7^2\ y_7\left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2\left(2 - n_7 + n_7 y_7\right)\right]$$

$$-W_7\ X_7\left[\frac{X_7(1 + n_7 y_7)}{M_7^2(2 + n_7 y_7)^2} + 2\ y_7\left(n_7 - 1\right) + \frac{X_7^2(7 + y_7)}{M_7^2(2 + n_7 y_7)} = 0.\ (A2.12)\right.$$

Since $n_7$ and $A_7$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7), the\ condition (A2.77?) can\ be\ written\ as$

$[\frac{X_7(E_7-7)}{E_7^2}]^2[\frac{X_7^2}{M_7^2} - 4\ y_7\ E_7^2\frac{X_7}{M_7^2} + 4y_7(n_7 - 2)E_7^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$2\ E_7[y_7 E_7 + (y_7\ E_7(7 + y_7)(E_7 - n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - y_7 E_7(7 + y_7)(E_7 - n_7))^{7/2}]$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0 $W_7$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7-7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2-2y_7\frac{(E_7+n_7^2y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right)-4y_7(E_7-n_7)$$

And has the solution

$$\frac{X_7}{M_7^2}>\frac{E_7}{(E_7-7)}[y_7(E_7+n_7^2y_7)+[y_7^2(E_7+n_7^2y_7)^2+4y_7(E_7-7)(E_7-n_7)]^{7/2}]$$

Given $E_7>n_7$this lower limit to $X_7/M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7/M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$in the appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7<n_7$ $(ie.,2A_7M_7^2<X_7-Y_7)$then (A2.7)and (A2.8)show that $Q_7$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2}>2E_7[y_7E_7-(y_7E_7(7+y_7)(E_7-n_7)]^{7/2}]$$

And

$$E_7>n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2 > 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.

If $A_7^2 < 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.+ $(Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$ and $X_7$ and $Y_7$ by $W_7^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7 \lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]\, e^{2T_7\lambda_7} + M_7 V_7\, e^{2T_7\lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i} Z_7$ *and* separate

(A3.1)into real and imaginary parts

$$Z_7^2 + (M_7 + V_7) Z_7 \sin(Z_7 T_7) - (U_7 M_7 + Y_7) + (M_7(U_7 + V_7) - X_7)$$

$$\mathrm{X} \cos(Z_7 T_7) - 2M_7 V_7 \cos^2(Z_7 T_7) + M_7 V_7 = 0$$

$$(M_7 + U_7) Z_7 - (M_7 + V_7) Z_7 \cos(Z_7 T_7) + (M_7(U_7 + V_7) - X)$$

$$\mathrm{X} \sin((Z_7 T_7) - M_7 V_7 \sin(Z_7 T_7) = 0$$

Using (A.3.3) to eliminate $Z_7$and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - cos\, Z_7\ T_7$and $A_7 = U_7 - V_7$we obtain

$$4\,M_7^2 V_7^2 W_7^4 + 2\,M_7 V_7\ W_7^3[(M_7 - V_7)^2 + A_7(M_7 + V_7) + 2(A_7 M_7 - X_7)]$$

$$+ W_7^2[M_7 - V_7)^2(2\,A_7\ M_7 + Y_7 - X_7) + (2A_7\ M_7 - X_7)^2 - A_7(M_7 - V_7)(2A_7\ M_7 - X_7)$$

$$+ M_7 V_7\,(A_7^2 + 4(X_7 + Y_7))]$$

$$+ W_7[A_7^2 + 2X_7)(2\,A_7\ M_7 - X_7) + 2A_7 Y_7(M_7 + V_7) - M_7 A_7^3] + A_7^2(X_7\ M_7 - Y_7) = 0$$

Being a quartic in $W_7$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7 > 0$$

$$and$$

$$C_7 > (X_7 - Y_7)/2\,M_7$$

$recalling\ that\ A_7 = U_7$and defining $C_7 = U_7 + V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0, U_7 V_7 = A_7^2(X_7 + Y_7)$$

$$Q_7\left(2, U_7 V_7\right) = \left(C_7 + 2M_7^2\right)^2 (2C_7 M_7 + Y_7 - X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

## Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

**Module numbered five**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This is the condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference

between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \quad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t - T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\left(R(t)\right)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

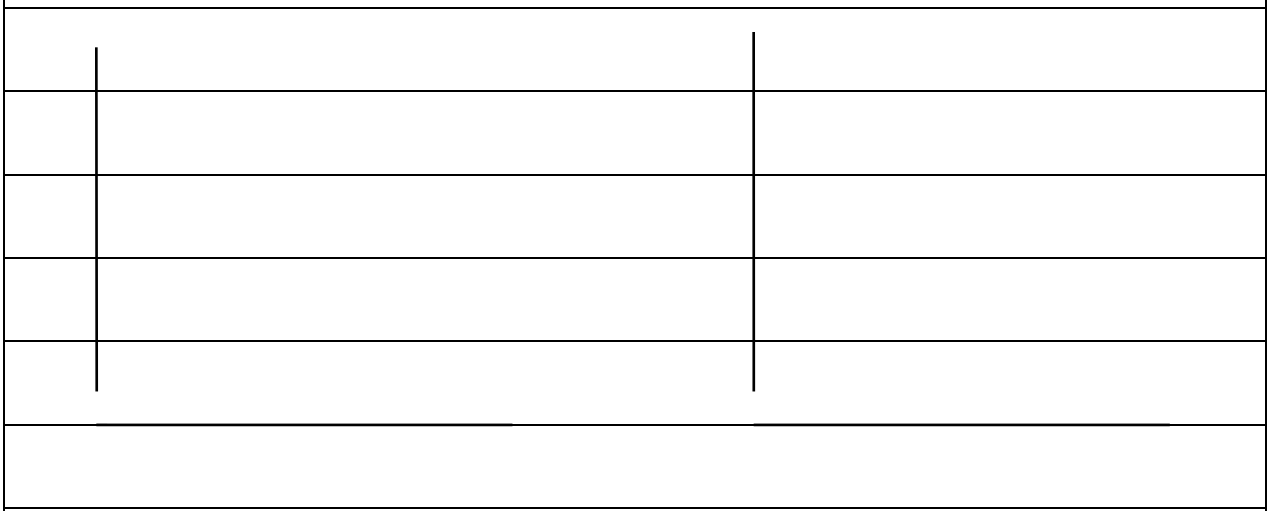
Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time

delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the



TIME DELAY

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period,

which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

As it is usually expressed outside of phenomenological texts, to undergo a conscious experience necessarily means that there is something it is like for the subject to have that experience (Nagel 1974; Searle 1992). This is obviously true of bodily sensations like pain. But it is also the case for perceptual experiences, experiences of desiring, feeling, and thinking. There is something it is like to taste chocolate, and this is different from what it is like to remember what it is like to taste chocolate, or to smell vanilla, to run, to stand still, to feel envious, nervous, depressed or happy, or to entertain an abstract belief. Yet, at the same time, as I live through these differences, there is something experiential that is, in some sense, the same, namely, their distinct first-personal character. All the experiences are characterized by a quality of mineness or for-me-ness, the fact that it is I who am having these experiences. All the experiences are given (at least tacitly) as my experiences, as experiences I am undergoing or living through. All of

this suggests that first-person experience presents me with an immediate and non-observational access to myself, and that consequently (phenomenal) consciousness consequently entails a (minimal) form of self-consciousness. To put it differently, unless a mental process is pre-reflectively self-conscious there will be nothing it is like to undergo the process, and it therefore cannot be a phenomenally conscious process. The mineness in question is not a quality like being scarlet, sour or soft. It doesn't refer to a specific experiential content, to a specific what; nor does it refer to the diachronic or synchronic sum of such content, or to some other relation that might obtain between the contents in question. Rather, it refers to the distinct givenness or the how it feels of experience. It refers to the first-personal presence or character of experience. It refers to the fact that the experiences I am living through are given differently (but not necessarily better) to me than to anybody else. It could consequently be claimed that anybody who denies the for-me-ness of experience simply fails to recognize an essential constitutive aspect of experience. Such a denial would be tantamount to a denial of the first-person perspective. It would entail the view that my own mind is either not given to me at all — I would be mind- or self-blind — or is presented to me in exactly the same way as the minds of others. There are also lines of argumentation in contemporary analytical philosophy of mind that is close to and consistent with the phenomenological conception of pre-reflective self-awareness. Alvin Goldman provides an example: [Consider] the case of thinking about x or attending to x. In the process of thinking about x there is already an implicit awareness that one is thinking about x. There is no need for reflection here, for taking a step back from thinking about x in order to examine it…When we are thinking about x, the mind is focused on x, not on our thinking of x. Nevertheless, the process of thinking about x carries with it a non-reflective self-awareness (Goldman 1970, 96). A similar view has been defended by Owen Flanagan, who not only argues that consciousness involves self-consciousness in the weak sense that there is something it is like for the subject to have the experience, but also speaks of the low-level self-consciousness involved in experiencing my experiences as mine (Flanagan 1992, 194). As Flanagan quite correctly points out, this primary type of self-consciousness should not be confused with the

much stronger notion of self-consciousness that is in play when we are thinking about our own narrative self. The latter form of reflective self-consciousness presupposes both conceptual knowledge and narrative competence. It requires maturation and socialization, and the ability to access and issue reports about the states, traits, dispositions that make one the person one is. To claim that every kind of self-consciousness is conceptual is overly cognitive. Bermúdez (1998), to mention one further philosopher in the analytic tradition, argues that there are a variety of **nonconceptual forms of self-consciousness** that are "logically and ontogenetically more primitive than the higher forms of self-consciousness that are usually the focus of philosophical debate" (1998, 274; also see Poellner 2003). This growing consensus across philosophical studies supports the phenomenological view of pre-reflective self-consciousness. (References: Stanford Encyclopedia, Post Modersnist Theories, Internet Encyclopedia of Philosophy) Korzybski pointed out that there are two ways to slice easily through life; to believe everything or to doubt everything. Both ways save us from thinking. So, from time to time you have to give yourself a Virtual Reality Check. The hallmark of postmodern philosophy has been disbelief or skepticism of all "metanarratives," or translations of reality. Postmodernism has even turned its profound skepticism on such important humanist concepts as "objective truth" and reason.Yet, for a deconstructionist postmodern society, individually we are still riddled with superstition and gullibility, and open to manipulation through our belief systems as any politician, philosopher, clergy, or salesperson will attest. Further, most people are painfully naive when it comes to even the simplest scientific understanding. Most of us don't have a clue about the fundamental nature of physical reality or our own psychological nature, and our ability to be fooled by our senses and mind, or methods of social and technological persuasion. We may think we know, but the penetration is deeper than we dare imagine. Nature rejects the naiveté that seeks absolute truth. We are beginning to realize, individually and culturally, that "realities" are all human constructions. The task becomes one of "catching ourselves in the act" of creating our own "reality" from the flow of events. Human truth is always an engagement of mind with experience.. We don't need to fear the collapse of our

personalistic belief system (the "box" we live in), nor our belief in absolute truth. **A strong desire to engage in the "quest for uncertainty" complements our anxiety that perhaps there is no absolute, objective ground to reality**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t-T_8)F_8(R_8(t-T_8) - N_8(t)M_8((R_8(t) \qquad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8} + \frac{N_8 dG_8}{dR_8} - \frac{dB_8}{dR_8} > \frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8} + \frac{dM_8}{dR_8}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_8 = \frac{N_8 G_8}{2M_8}\frac{d(8nF_8)}{dR_8} (\geq 0) \quad (8a)$$

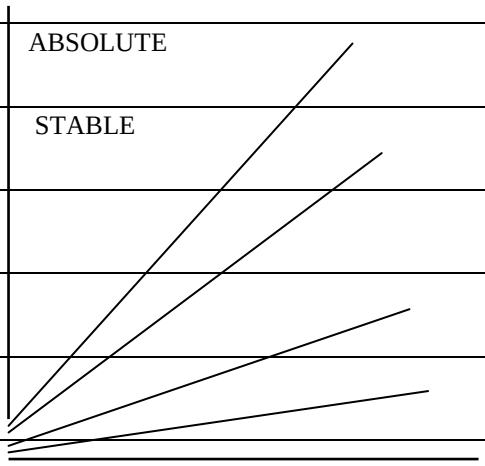
$$I_8 = -\frac{N_8 G_8}{2M_8}\frac{d(8nM_8}{dR_8} (\geq 0) \quad (8b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8 > L_8 - I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_8 > L_8 - I_8$, when $P_8^2/4 > I$ (8a)

$P_8 > L_8 - I_8 + \frac{(I_8 - P_8^2}{(L_8 - I)}$ when $P_8^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

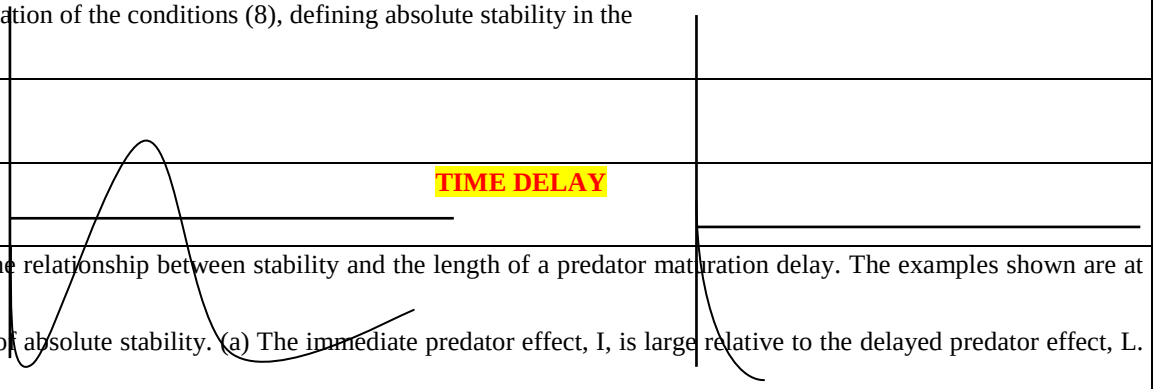



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this

rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior , we must note that the model could be applied to psychological systems too, albeit we are still not sure of the fact how the wave function collapses thereby due to perception.

Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility. He privatizes the sublime, having learned from the fate of his predecessors that the public can never be more than beautiful."(11) Notice is made here of what is indeed the telling feature. Levinas is praised, for example, for his alarm when "the social will [is] sought in an ideal of fusion . . . the subject losing himself in a collective representation, in a common ideal. . . . It is the collectivity which says 'us,' and which, turned toward the intelligible sun, toward the truth, experience, the other at his side and not face to face with him. The tellingly titled D'un ton apocalyptique contains passages that are particularly excited: These people situate themselves outside the ordinary, but they have in common this: they describe themselves as having an immediate and intuitive relation with mystery. And they want to attract, to seduce, and lead others to the mystery, through mystery. This agogic function of the leader of men, of the duce, of the Führer, of the leader, places him above the crowd that he manipulates with the aid of a small numbers of adepts joined together in a sect with a secret language, a clique or a small party with its

ritualized practices. The mystifiers pretend to have exclusive access to the privilege of a secret mystery. . . . The revelation or the unveiling of the secret is something that they jealously reserve for themselves. Jealousy is here a major characteristic. There is not narcissism and non-narcissism; there are narcissisms that are more or less comprehensive, generous, open, extended. What is called non-narcissism is in general but the economy of a much more welcoming, hospitable narcissism, one that is much more open to the experience of the other as other. I believe that without a movement of narcissistic reappropriation, the relation to the other would be absolutely destroyed; it would be destroyed in advance. The relation to the other--even if it remains asymmetrical, open, without reapporpriation--must trace a movement of reappropriation in the image of oneself for love to be possible, for example. Love is narcissistic. Beyond that, there are little narcissisms, there are big narcissisms. . . .(23) **Justice of the Pieces Deconstruction as a Social Psychology Douglas Collins** Sabine Wilke cogently argues for the rise of the aesthetic dimension in philosophical writing by showing how the metaphorical and poetic style of Heidegger, the principles of phrase configurations with Adorno, or the topographic writing of derrida replace through their very form something what traditionally would have been subject to discursive exposure The concern shifts from telling to showing, showing first and foremost that something remains untold. An exemplary book in this aesthetic drive is Avital Ronell's Telephone Book where the playful layout and explicit break with traditional design interrupt the sequential logic of the paragraphs thereby performing the very argument at stake: open up new circuits of meaning and signification. The engagement with form and publishing formats is only consequential since media theory as exposed by Walter J. Ong, Eric A. Havelock, Parry Miman or Marshall McLuhan besides many others made it clear that different communication technologies condition the mode of thought of a given culture The fractals of Benoit Mandelbrot or the butterfly shaped Lorenz attractor are the popular icons of this turn in the sciences. For communication studies and philosophy, Vilém Flusser was one of the pioneer thinkers reflecting the rise of what he called “techno-images.” The photograph is the paradigmatic case which condenses the whole crises of Western

thought and culture on its grainy surface. It subverts the linear logic of writing.6 Flusser's central diagnosis detects a shift from the grand narratives associated with the literate mind and Modernity towards creating and envisioning the world through image based practices. Writing himself in a succinct, visually inspired style, Flusser contents that the web of meaning, spun throughout the centuries by our textual practices, has become transparent. We discover the void between the letters through which the world escapes: "We are alienated from the world of text, precisely because we see through them as our own product."7 As a consequence we are confronted with a world of dissociated facts and information bits. Flusser hails this fragmented world as the advent of the zero dimensional point-universe that will provide the new raw material for our image based practices. Techno-images are nothing else than a conscious, inherently artistic drawing-together of point-elements as exemplified first and foremost by the silver-grains of the photograph but also by more recent pixel based techno-images. Similar to blueprints, designs or statistical curves, the weather map for example visualizes and contracts a tremendous amount of meteorological data from widely distributed stations into a single image of high- and low pressure

areas, temperatures, wind speeds, etc. Techno-images "compute" in this sense the universe of unimaginable, abstract information-bits into a concrete surface – the image. Echoing Paul Feyerabend, these image practices have to be understood as creative practices where the truth of scientific models has been replaced by its aesthetic quality **Aesthetics, Writing, Networked Computers Jörg Muller Dissertation Submitted to the Division of Media and Communications of The European Graduate School in Candidacy for the Degree of Doctor of Philosophy**

**Lothar Schäfer (1) and Sisir Roy (2) in** Quantum Reality, the Importance of Consciousness in the Universe, the Discovery of a Non-Empirical Realm of Physical Reality, and the Convergence with Ancient Traditions of Indian and Western Philosophy investigate metaphysical aspects that include 1) the discovery of a non-empirical part of physical reality in a realm of potentiality; 2) the emanation of the empirical world out of a realm of non-material forms; 3) the discovery that the nature of physical reality is that of an indivisible Wholeness – the One; and 4)

panpsychism: the possibility that the One is aware of its processes like a Cosmic Consciousness. The convergence of powerful traditions of seemingly disparate cultures is particularly important to point out in the present process of globalization, when a unifying view is needed to avoid controversy and conflict. To quote them in extenso : This is the nonlocality of the quantum world. Menas Kafatos and Robert Nadeau (1990) have drawn a remarkable conclusion from this phenomenon: If reality is nonlocal, the nature of the universe is that of an undivided wholeness. Because our consciousness has emerged from this wholeness and is part of it, it is possible to conclude that an element of consciousness is active in the universe: a Cosmic Consciousness. Pre-reflective self-consciousness is pre-reflective in the sense that (1) it is an awareness we have before we do any reflecting on our experience; (2) it is an implicit and first-order awareness rather than an explicit or higher-order form of self-consciousness. Indeed, an explicit reflective self-consciousness is possible only because there is a pre-reflective self-awareness that is an on-going and more primary self-consciousness. Although phenomenologists do not always agree on important questions about method, focus, or even whether there is an ego or self, they are in close to unanimous agreement about the idea that the experiential dimension always involves such an implicit pre-reflective self-awareness. In line with Edmund Husserl (1959, 189, 412), who maintains that consciousness always involves a self-appearance (Für-sich-selbst-erscheinens), and in agreement with Michel Henry (1963, 1965), who notes that experience is always self-manifesting, and with Maurice Merleau-Ponty who states that consciousness is always given to itself and that the word ‘consciousness’ has no meaning independently of this self-givenness (Merleau-Ponty 1945, 488), Jean-Paul Sartre writes that pre-reflective self-consciousness is not simply a quality added to the experience, an accessory; rather, it constitutes the very mode of being of the experience: This self-consciousness we ought to consider not as a new consciousness, but as the only mode of existence which is possible for a consciousness of something (Sartre 1943, 20 [1956, liv]). The notion of pre-reflective self-awareness is related to the idea that

experiences have a subjective 'feel' to them, a certain (phenomenal) quality of 'what it is like' or what it 'feels' like to have them.

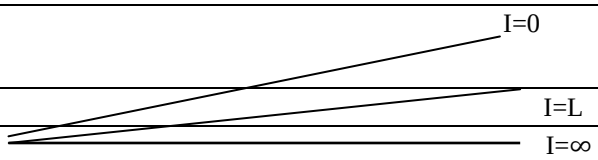
Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

TOTAL

ABSOLUTELY STABLE

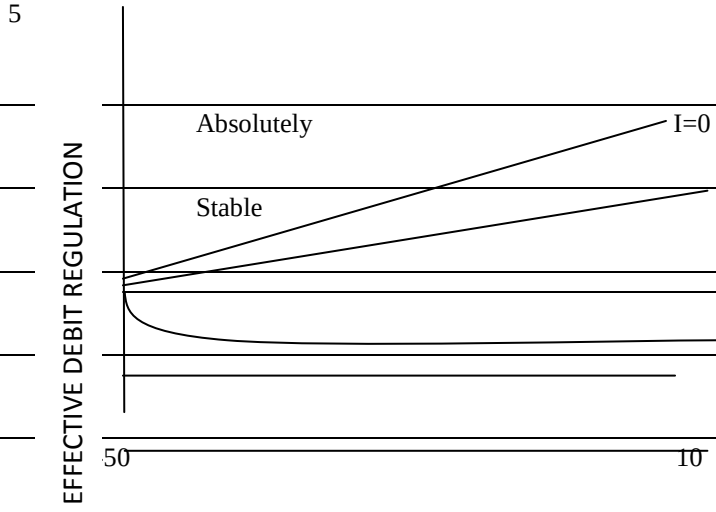



**DELAYED DEBIT EFFECT**

FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

## Utility Analysis And neural Network

Caroline Williams investigated affective density of the political and its effect on our understanding of political subjectivity. Taking up Spinoza's challenge to think about affect beyond corporeal embodiment, he argues that there is a modality of affectivity that cannot simply be inscribed within the borders of subjectivity. theorising affect as an impersonal force anchored in a relational ontology that gives due recognition to the circulation of affects, as well as to their ambivalent structure in creating sites of identification, and he utilises this ontology to reflect on the dynamic of the political and the shape of political subjectivity. Williams argues that Spinoza's philosophy (through ideas of conatus and imagination) offers the conceptual resources to reconfigure the composition of affective subjectivity as a transindividual social bond and as an unconscious dynamic of ethico-political existence. **Subjectivity (2010) 3, 245–262. doi:10.1057/sub.2010.15 Affective processes without a subject: Rethinking the relation between subjectivity and affect with Spinoza Caroline Williams** Phenomenological Time bears ample evidence and infallible observatory and apodictic evidence albeit with a sense of chagrined circumvention and defeated disgruntlment circularity of both psychological and phenomenological interpretations of art which are sometimes insightful, but most often work on the logic of cyclical time and unavoidable drives and representations played out as a ‘natural’ order: where ‘men can’t avoid being boys’ a circularity which valorizes a helplessness/truthfulness conundrum. One begins with a classical phenomenological fallacy with Duchamp’s Fountain, does the urinal keep returning or do we keep returning to the urinal? Put in this way, it is a male question, something that phenomenology conveniently conceals when positing the “I”, or indeed, the “we”. But we must be careful not to simply assume that the phenomenological first person is white, heterosexual, able-bodied and male, providing a foundation upon which knowledge should be based as a universal condition. Deleuze’s second critique of time is linked to Kant. This model of time extends not only to phenomenology but to the phenomenological interpretation

of art and is seen as a straight line. In the Critique of Pure Reason, Kant cuts up circular time and reassembles it as a series of sensory experiences which memory processes and which constitute subjectivity. Whereas the circular model of time (and interpretation) inscribes fatalism into the subject and the world, the phenomenological method is in danger of solipsism, where the subject constitutes the world. Much of embodied philosophy and enactivist theories of cognition go down this line with a plethora of art historical analyses by the likes of Rosalind Krauss, Michael Fried and others taking the lead from Merleau-Ponty. These approaches successfully challenge the **homunculus model of a Cartesian mind-body** dualism where there is a supervisor in the mind who processes experience and even consciousness and to whom all representations are directed. But phenomenology is mainly concerned with perceptual consciousness, the processing of colors, shapes, forms, motion in art and the world. The senses and sensorimotor processes are the fundamental pathways by which the interior is connected to the exterior in a chiasm of self-constitution in the world. The problem with this is that we explore the art through our sense perceptions, and what does art reveal to us?—our sense perceptions! Such solipsistic approaches have real problems relating complex and abstract conceptual production to bodily and sensorimotor contingencies, and they nearly always posit the subject as the site and object of discovery in the world. To be fair, phenomenology has moved on to rather more sophisticated models of intersubjectivity, taking into account the specificities of different kinds of bodies and ways of being beyond the white, male heterosexual as the common denominator, yet many naïve approaches to art perpetuate the notion that perceptual, rather than conceptual processes and evolutionary neural patterns are primary, fundamental ways to understand the art experience In fact, these aspects are only part of art production and reception, aspects which are often filtered out or made known to us in order for us to parody such trigger responses, as in much of Duchamp's work, which holds in contempt such retinal approaches to art. If we take the **subject of abstract expressionism** in the 1950s or later minimalist art, art history, teaching and interpretation still insists on putting the act of perceiving as fundamental to the constitution not only of the phenomenological first

person but also of art itself. The male act of usually standing up, which consists of directing a stream with the hand(s) into a container or on a flat surface, sand or snow in order to leave patterns and traces (the usual child's prank) underlies the logic of macho 1950s Abstract Expressionist culture, especially that of Jackson Pollock, so effectively parodied by Carolee Schneeman's Vagina Scroll, 1975. Here, painting or writing is a kind of a discharge which invests the **object and space with** the phenomenological "I", disciplined by existential struggle and reflected back by the traces, signatures and gestures left behind in the work of art, rediscoverable as an element in the making of other subjectivities.It is precisely this kind of solipsism which is blind to the art event as something well in excess of re-establishing habits of subjective self-identity through the production and interpretation of art. The attempt to reinforce a phenomenological identity upon the event by personalizing sense experience is unable to see art as something which remodels the phenomenological 'I' and neural plasticity itself. We learn new technologies and redraw our neural pathways through such new and spontaneous experiences which draw us into intersubjective worlds. Duchamp's Fountain and Deleuze's Repetition and Différence Grégory Minis sale (I have made some changes) The relation of different groups to deconstruction is settled by its relation to groups. An American enthusiast--one who has had a role in the development of our affection for Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^{\infty} \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$then given the constraints of $f(T_8)$, $(A1.5) can\ never\ be\ satisfied\ Co$nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8 \lambda_8 e^{-T_8 \lambda_8} + [(U_8 - V_8) M_8 + Y_8] - [((U_8 - V_8) M_8 -] e^{-T_8 \lambda_8} = 0$$

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$\boldsymbol{Z}_8^2 + M_8 Z_8 \sin(Z_8 T_8) - (A_8 M_8 + Y_8) + (A_8 M_8 + X_8) \cos(Z_8 T_8) = 0$$

$$(A_8 M_8) Z_8 - M_8 Z_8 \cos(Z_8 T_8) + (A_8 M_8 - X_8) \sin(Z_8 T_8) = 0$$

Where $A_8 = U_8 - V_8$ Using (A2.3), substitute for $\boldsymbol{Z}_8^2$ and $Z_8$ in (A2.2):

$W_8^2 (X_8^2 - 3 M_8 A_8 X_8 - M_8^2 (X_8 - Y_8) + 2 M_8^3 A_8 + 2 A_8^2 M_8^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_8 = 1 - \cos(Z_8 T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8 / 2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_8 \le 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2 (X_8 - Y_8)$$

$$Q_8(2, A_8) = (A_8 + 2M_8)^2 (2A_8 M_8 - (X_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2 \left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8 \left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8 \left(2X_8 + Y_8\right)\right)$$

$$+2A_8 \left(X_8 + Y_8\right)$$

$$= 2M_8^2 W_8^2 (M_8 + A_8) + M_8 A_8^2 W_8 + (2A_8 M_8 - X_8 + Y_8)(A_8 W_8 + M_8 W_8^2$$

$$+ (2 - W_8)((2M_8 X_8 + M_8 Y_8) W_8 + A_8 (X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $W_8$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in *A* will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range $0 \le W_8 \le 2$.Defining

$Q_8(W_8, A_8) = a_8 W_8^2 + c_8$ *where* $a_8$ $b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \square - b_8 \square 0$$

And

$$b_8^2 \square 4a_8 c_8$$

$$Case(i): Y_8 = 0$$

$Let\ Y_8 = 0$ And try the substitution A= $X_8 = /2M_8\, in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8 /$ $2M_8$) is positive for *W<2. Since* $Q_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\ X_8$ ,

$Let\ Y_8 = \ y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 + \ n_8 y_8)$ we have

$$W_8^2\, y_8 \left[X_8 n_8 \frac{(1 + n_8 y_8)}{(2 + n_8 y_8)} + M_8^2 \left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8 X_8 \left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2}+2y_8\left(n_8-1\right)+\frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)}\right]=0.\ (A2.12)$$

Since $n_8$ and $A_8$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8=(2+n_8 y_8)$, *the condition*$(A2.88?)$*can be written as*

$$\left[\frac{X_8(E_8-8)}{E_8^2}\right]^2\left[\frac{X_8^2}{M_8^2}-4y_8 E_8^2\frac{X_8}{M_8^2}+4y_8(n_8-2)E_8^3\right]$$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8/W_8^2)$ can be expressed as

$$2E_8\left[y_8E_8+(y_8E_8(8+y_8)(E_8-n_8))^{8/2}\right]>\frac{X_8}{M_8^2}>2E_8\left[y_8E_8-y_8E_8(8+y_8)(E_8-n_8))^{8/2}\right]$$

Thoughts on the Social and Political Implications of Correlationism .Lion Mirror Truth be told, as my thought has evolved the issue of correlationism had fallen off the radar for me. Somehow the debate had come to seem too "philosophical" to me, too "scholastic", too remote from what interests me: understanding why social assemblages are organized as they are, how power functions in social assemblages, and what we might do to address that power and change things. Somehow the question of whether or not we can get out of the correlation between thinking and being just came to seem remote from these sorts of issues. Somehow it seemed too epistemological. ASIDE: Numerous discussions over the years have led me to believe that the debate over correlationism is poorly understood (or maybe I just don't understand it). On countless occasions I've heard people say "of course we must relate to things in order to know them." Well yeah, of course! I don't think this is what the critic of correlationism is getting at. It seems to me that correlationism is something more robust than the theses that we must relate to something to know it.

Correlationism instead seems to require the theses that thought and being are indiscernible. Put more concretely, the correlationist is someone who argues that we either a) can never tell whether being is merely a construction of our thought (weak correlationism), or b) who argues that thought actually constructs being (strong correlationism). In other words, correlationism is another name for idealism. One can hold that we must relate to something in order to know it without being a correlationist. As an aside I should also add that I am a correlationist about some things. For example, I think money is something constructed by society and is therefore a strong correlationist when it comes to money. At any rate, for a long time I'd become rather indifferent to debates about correlationism and philosophies of access. I had learned the lessons of speculative realism– which I could have also learned, I think, from Deleuze and Guattari, the new materialist feminists, actor-network theorists such as Latour –and had moved on. However, occasionally you come across a tone of phrase that pitches something in a different light. In The Cut of the Real, Katerina Kolozova writes, …the political problem of contemporary philosophy identified by the 'new realists' is, in fact, the product of a more fundamental epistemic problem. In his book After Finitude, Quentin Meillassoux calls this problem 'correlationism' and identifies it as an essentially post-Kantian legacy, which continues to dominate and limit philosophy. As a matter of fact, correlationism lies at (eb) the heart of postmodern theory and consists in the premise that thought can only 'think itself,' that the real is inaccessible to knowledge and human subjectivity, and that there is nothing but discursive constructs that fully determine thinking and that are meth0dologically accounted for all the way down. (1 – 2) Thought thinking only itself: Thought only encountering itself. In the jargon of postmodern and poststructuralist lingo, this would be the thesis of infinite semiosis, where signs ("thoughts") only ever relate to other signs. Within this framework, discursivity comes to be the hegemonic framework defining all of being. At the level of politics and social theory more generally, if the correlationist thesis is true the

consequences are clear: all social phenomena are discursive and all solutions to social and political problems will be discursive. The sole sphere of the political will be the discursive and all questions of politics will be questions of speech-acts and interpretation. The problem here is not that many theorists recognize that the discursive and semiotic plays an important role in the social and the political. It does and I've repeated this tirelessly. The problem is with what happens when thought or the semiotic becomes a hegemon, an "all", foreclosing our ability to recognize other forms of power. What I've wanted to say is that not all power functions discursively. In my last post and elsewhere I spoke of some other forms of politics: Thermopolitics: The politics surrounding **energy in the form of calories and fuels such as gasoline** and coal, and how our life and our very bodies are structured by energy dependencies and by being trapped in particular distributive networks that render these forms of energy available. I'm being quite literal when I speak of energy, talking about the effects, for example, of the absence of food in certain educational environments on cognition, for example; and am generally hostile to metaphorical extensions of the concept of energy which I see as erasing the dimension of real materiality. Geopolitics: The role that features of natural and built geography such as mountain ranges, rivers, oceans, soil conditions, roads, housing design, etc., play in the form that social relations take and how they impact individual bodies. Chronopolitics: The way in which the structurations of time organize what is possible for us. For example, the structuration of the working day, how much we can say and comprehend at any given time, the impact of things like the invention of the clock, etc. **Oikopolitics: This would be the domain of political economy described so well by Marxists**. So five different types of politics: Semiopolitics (or what currently dominates critical theory), thermopolitics, geopolitics, chronopolitics, and oikopolitics. No doubt there are other sites of the political or political struggle that we could speak of, but this is a good start. Also, it should be obvious that these aren't exclusive domains, but are entangled in all sorts of important

ways. For example, something might take place at the level of semiopolitics (speech, law, rhetoric, norms, and communication) that has all sorts of impact at the level of thermopolitics. Congress might decide to cut programs that fund school meal programs. This, in turn, will have a thermodynamic impact on those students that go without the calories they need developmentally and cognitively to function in a particular way. There is an entanglement here of semiopolitical and thermopolitical domains. The young student here has been constrained both at the level of **semiotic phenomena and thermodynamic structures**. The point is that if true, semiotic intervention (speech-acts, protests, interpretations, deconstructions, etc) will not be an appropriate response to all political problems because social formations are not entirely structured by the semiotic. The child in that school does not suffer from a lack of the right signs, but from a lack of calories needed to run the engine of his thought and body. Certainly semiotic interventions might be needed to render that energy available, but it is the energy itself that is at issue and the absence of that energy that forms the spider web entangling him in his position. A correlationist perspective tends to erase this as even being a site of the political. **Larval Subjects January 28, 2014 Ontological Anarché Posted by larval subjects under Uncategorized January 27, 2014 Roden on Pluralism Over at enemy industry Roden has an excellent post up on the pluralism discussion. Check it out here. January 25, 2014 Different Senses of Pluralism and Ontology Posted by larval subjects under Uncategorized Responding to one of my comments over at his blog, Jeremy Trombley:** I have deleted restricted, reformulated some sentences .Kindly bear with me. Intention is to make the discursive enucleation consistent with the subject in question**.** Recourse to epistemology of the quantum mechanics generic and fervor of string theory is a sine qua non for the further comprehension of the theory and ontology .How would the world appear to us if its ontology was that of classical mechanics but every agent faced a restriction on how much they could come to know about the classical state? We show that in most respects

it would appear to us as quantum. The statistical theory of classical mechanics, which specifies how probability distributions over **phase space evolve** under Hamiltonian evolution and under measurements, is typically called Liouville mechanics, so the theory we explore here is Liouville mechanics with an epistemic restriction. The particular epistemic restriction we posit as our foundational postulate specifies two constraints. The first constraint is a classical analog of Heisenberg's uncertainty principle; the second-order moments of position and momentum defined by the phase-space distribution that characterizes an agent's knowledge are required to satisfy the same constraints as are satisfied by the moments of position and momentum observables for a quantum state. The second constraint is that the distribution should have maximal entropy for the given moments. Starting from this postulate, authors derive the allowed preparations, measurements, and transformations and demonstrate that they are isomorphic to those allowed in Gaussian quantum mechanics and generate the same experimental statistics. They argue that this reconstruction of Gaussian quantum mechanics constitutes additional evidence in favor of a research program wherein quantum states are interpreted as states of incomplete knowledge and that the phenomena that do not arise in Gaussian quantum mechanics provide the best clues for how one might reconstruct the full quantum theory. I: http://dx.doi.org/10.1103/PhysRevA.86.012103 **Reconstruction of Gaussian quantum mechanics from Liouville mechanics with an epistemic restriction Phys. Rev. A 86, 012103 – Published 10 July 2012 Stephen D. Bartlett, Terry Rudolph, and Robert W. Spekkens.** Deleuze in fact characterizes modes of existence, with their powers and capacities. The answer is this: Deleuze approaches modes of existence, ethically speaking, not in terms of their will, or their conscious decision making power (as in Kant), nor in terms of their interests (as in Marx, for example), but rather in terms of their drives. For Deleuze, conscious will and preconscious interest are both subsequent to our unconscious drives, and it is at the level of the drives that we have to aim our ethical analysis. Here, I would like to

focus on two sets of texts on the drives taken, not from Nietzsche and Spinoza, but rather from Nietzsche and Leibniz (Leibniz being one of the first philosophers in the history of philosophy to have developed a theory of the unconscious). The first set of texts comes from Nietzsche's great early book entitled Daybreak, published in July 1881. Nietzsche first approaches the question of the drives by giving us an everyday scenario: "Suppose we were in the market place one day," he writes, "and we noticed someone laughing at us as we went by: this event will signify this or that to us according to whether this or that drive happens at that moment to be at its height in us—and it will be a quite different event according to the kind of person we are. One person will absorb it like a drop of rain, another will shake it from him like an insect, another will try to pick a quarrel, another will examine his clothing to see if there is anything about it that might give rise to laughter, another will be led to reflect on the nature of laughter as such, another will be glad to have involuntarily augmented the amount of cheerfulness and sunshine in the world—and in each case, a drive has gratified itself, whether it be the drive to annoyance, or to combativeness or to reflection or to benevolence. This drive seized the event as its prey. Why precisely this one? Because, thirsty and hungry, it was lying in wait" (D 119). This is the source of Nietzsche's doctrine of perspectivism ("there are no facts, only interpretations"), but what is often overlooked is that, for Nietzsche, it is our drives that interpret the world, that are perspectival—and not our egos, not our conscious opinions. It is not so much that I have a different perspective on the world than you; it is rather that each of us has multiple perspectives on the world because of the multiplicity of our drives—drives that are often contradictory among themselves. "Within ourselves," Nietzsche writes, "we can be egoistic or altruistic, hard-hearted, magnanimous, just, lenient, insincere, can cause pain or give pleasure" (Parkes, pp. 291-292). We all contain such "a vast confusion of contradictory drives" (WP 259) that we are, as Nietzsche liked to say, multiplicities, and not unities. Moreover, these drives are in a constant struggle or combat with each other:

my drive to smoke and get my nicotine rush is in combat with (but also coexistent with) my drive to quit. This is where Nietzsche first developed his concept of the will to power—at the level of the drives. "Every drive is a kind of lust to rule," he writes, "each one has its perspective that it would like to compel all the other drives to accept as a norm" (WP 481). DELEUZE AND THE QUESTION OF DESIRE: TOWARD AN IMMANENT THEORY OF ETHICS Daniel W. Smith. **In on a Possible Physical Metatheory of Consciousness Miroljub Dugic et al** show that the modern quantum mechanics, and particularly the theory of decoherence, allows formulating a sort of a physical metatheory of consciousness. Particularly, the analysis of the necessary conditions for the occurrence of decoherence, along with the hypothesis that consciousness bears (more-or-less) well definable physical origin, leads to a wider physical picture naturally involving consciousness. This can be considered as a sort of a psycho-physical parallelism, but on very wide scales bearing some cosmological relevance. Open Systems & Information Dynamics Volume 9, Number 2, 153 (2002) arXiv: quant-ph/0212128. On a broader perspective and background, the theoretical computation of the universe is justified.IN physics and cosmology, digital physics is a collection of theoretical perspectives based on the premise that the universe is, at heart, describable by information, and is therefore computable. Therefore, according to this theory, the universe can be conceived of as either the output of a deterministic or probabilistic computer program, a vast, digital computation device, or mathematically isomorphic to such a device. Digital physics is grounded in one or more of the following hypotheses; listed in order of decreasing strength. The universe or reality: is essentially informational (although not every informational ontology needs to be digital); is essentially computable (the pancomputationalist position); can be described digitally; is in essence digital; is itself a computer (pancomputationalism); is the output of a simulated reality exercise. Following Jaynes and Weizsäcker, the physicist John Archibald Wheeler wrote the following: [...] it is not unreasonable to imagine that

information sits at the core of physics, just as it sits at the core of a computer. (John Archibald Wheeler 1998: 340) It from bit. Otherwise put, every 'it'—every particle, every field of force, even the space-time continuum itself—derives its function, its meaning, its very existence entirely—even if in some contexts indirectly—from the apparatus-elicited answers to yes-or-no questions, binary choices, bits. 'It from bit' symbolizes the idea that every item of the physical world has at bottom—a very deep bottom, in most instances—an immaterial source and explanation; that which we call reality arises in the last analysis from the posing of yes–no questions and the registering of equipment-evoked responses; in short, that all things physical are information-theoretic in origin and that this is a participatory universe. (John Archibald Wheeler 1990: 5) David Chalmers of the Australian National University summarised Wheeler's views as follows: Wheeler (1990) has suggested that information is fundamental to the physics of the universe. According to this 'it from bit' doctrine, the laws of physics can be cast in terms of information, postulating different states that give rise to different effects without actually saying what those states are. It is only their position in an information space that counts. If so, then information is a natural candidate to also play a role in a fundamental theory of consciousness. We are led to a conception of the world on which information is truly fundamental, and on which it has two basic aspects, corresponding to the physical and the phenomenal features of the world. (Wikipedia). **Participatory consciousness forms the bastion, pillar, post, stylobate and fulcrum of the redressal measures and corrective steps taken for circumvention of many problems, identified and unidentified, spread widest commonalty. Anne Baring's sententious and pithy article will be good to have a timely outlook.** I would like to open my contribution to this meeting with the words of Ken Wilber that I quoted in the first Millennium Symposium: 'The secret impulse of life is towards greater consciousness. Maybe the evolutionary sequence is from matter to body, to mind, to soul, to spirit - each transcending and including, each with a greater depth and greater

consciousness and wider embrace' Ken Wilber, and with these from Rupert Sheldrake's contribution to the same meeting: "It is possible to make quite a good case that the sun could be conscious and if the sun, then why not the stars? Why not the galaxies? Why should there not be galactic minds? And if they communicate with each other then the distances are far too great for light to be the medium of transmission. We are into inter-galactic telepathy." After many decades of study, I know of only one diagram in all religious traditions that can adequately explain these statements; that can give us both a blue-print of the many levels of consciousness and a route-map to help us reach beyond our present level of understanding. I would like to introduce a diagram of the four worlds of Kabbalah to illustrate the theme that spirit has brought these four worlds into being and that we, living in the fourth and manifest world know nothing of the existence of the other three invisible ones. It is the clearest diagram I know of that can explain to us the multi-levelled structure of consciousness and our place in the great chain of being. We are indeed, as Wilbur suggests, evolving from matter to body, to mind, to soul, to spirit but where have we come from? It is possible that, as participants in the emanation of spirit into manifest life, we have also come from spirit to soul, to mind, to body, to matter. Everything, from the immense galactic energies to the minute forms of matter participates in one life, is generated from one creative spirit. There is nothing outside this life, this spirit. This diagram suggests that truth is not an article of faith but the experience of participating in the undiscovered reality of what has brought us into being. We cannot know truth until we have entered into communion with that reality. Continuing in similar vein **Michel Bitbol dilates upon the dialectic deliberation and polemical argumentation between** hermeneutists and eliminativists. When he formulated the program of Neurophenomenology, Francisco Varela suggested a balanced methodological dissolution of the “hard problem” of consciousness. He shows that his dissolution is a paradigm which imposes itself onto seemingly opposite views, including materialist approaches. I also point out that

Varela's revolutionary epistemological ideas are gaining wider acceptance as a side effect of a recent controversy between hermeneutists and eliminativists. Finally, he emphasizes a structural parallel between the science of consciousness and the distinctive features of quantum mechanics. This parallel, together with the former convergences, point towards the common origin of the main puzzles of both quantum mechanics and the philosophy of mind: neglect of the constitutive blindspot of objective knowledge. Phenomenology and the Cognitive Sciences Phenomenology and the Cognitive Sciences **Phenomenology and the Cognitive Sciences 2002, Volume 1, Issue 2, pp 181-224 Science as if situation mattered Michel Bitbol**

## ABSOLUTE STABILITY

Hence if $X_8 / M_8^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_8$ ) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8 - 8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2 - 2y_8\frac{(E_8 + n_8^2 y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right) - 4y_8(E_8 - n_8)$$

And has the solution

$$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8 - 8)}[y_8(E_8 + n_8^2 y_8) + [y_8^2 (E_8 + n_8^2 y_8)^2 + 4 y_8 (E_8 - 8)(E_8 - n_8)]^{8/2}]$$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2E_8[y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSOILIDATION**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + (M_8 + U_8)\lambda_8 - (M_8 + V_8)\lambda_8 e^{2T_8\lambda_8} + (M_8 U_8 + Y_8)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8\, e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + (M_8 + V_8) Z_8 \sin(Z_8 T_8) - (U_8 M_8 + Y_8) + (M_8(U_8 + V_8) - X_8)$$

$$\mathrm{X}\cos(Z_8 T_8) - 2M_8 V_8 \cos^2(Z_8 T_8) + M_8 V_8 = 0$$

$$(M_8 + U_8) Z_8 - (M_8 + V_8) Z_8 \cos(Z_8 T_8) + (M_8(U_8 + V_8) - X)$$

$$\mathrm{X}\sin((Z_8 T_8) - M_8 V_8 \sin(Z_8 T_8) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 =$

$1 - \cos Z_8\, T_8$and $A_8 = U_8 - V_8$we obtain

$$4M_8^2 V_8^2 W_8^4 + 2M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8(M_8 + V_8) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2[M_8 - V_8)^2(2A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8(M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8(A_8^2 + 4(X_8 + Y_8))]$$

$$+W_8\left[A_8^2+2X_8\right)\left(2A_8M_8-X_8\right)+2A_8Y_8\left(M_8+V_8\right)-M_8A_8^3]+A_8^2\left(X_8M_8-Y_8\right)=0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8>0$$

$$and$$

$$C_8>(X_8-Y_8)/2M_8$$

$recalling\ that\ A_8$= $U_8$and defining $C_8=U_8+V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8(0,U_8V_8=A_8^2(X_8+Y_8)$$

$$Q_8\left(2,U_8V_8\right)=\left(C_8+2M_8^2\right)^2(2C_8M_8+Y_8-X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0$W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$the existence of a D-position at this point is determined by letting cos $(Z_8T_8)=-1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all $T$, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**ABSOLUTE STABILITY IN PREDATOR- PREY MODELS**

**APPLIED TO ACCETUATION ATTRITION MODELS**

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability

criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The Percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending

upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t-T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\big(R(t)\big)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The

regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the

**TIME DELAY**

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a dominant eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

DELAY DUE TO CONGREGATIONAL AND CONFEDERATIONAL DEBIT CREDIT TYEPES

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t-T_8)F_8(R_8(t-T_8) - N_8(t)M_8((R_8(t) \quad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8} + \frac{N_8 dG_8}{dR_8} - \frac{dB_8}{dR_8} > \frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8} + \frac{dM_8}{dR_8}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_8 = \frac{N_8 G_8}{2M_8}\frac{d(8nF_8)}{dR_8} (\geq 0) \quad (8a)$$

$$I_8 = -\frac{N_8 G_8}{2M_8}\frac{d(8nM_8}{dR_8} (\geq 0) \quad (8b)$$

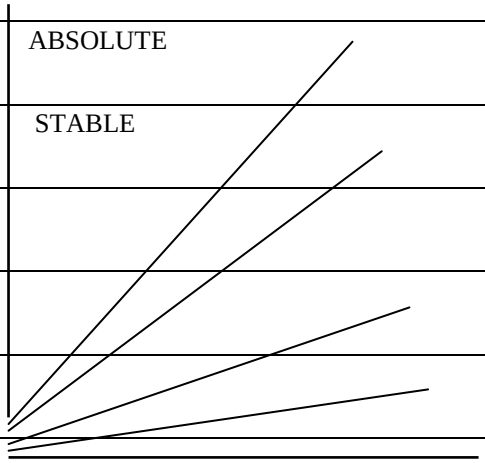
 

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8 > L_8 - I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_8 > L_8 - I_8,$ when $P_8^2/4 > I$ (8a)

$P_8 > L_8 - I_8 + \frac{(I_8 - P_8^2}{(L_8 - I)}$ when $P_8^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit

(like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

See the diagram Nunney gives for predator and prey analysis. We can simulate the same for the financial system in the foregoing too.

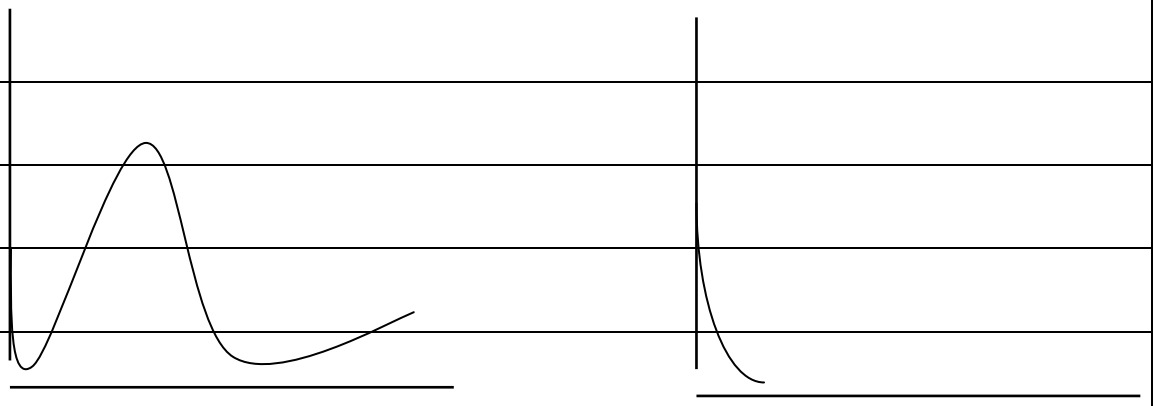
Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the

debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

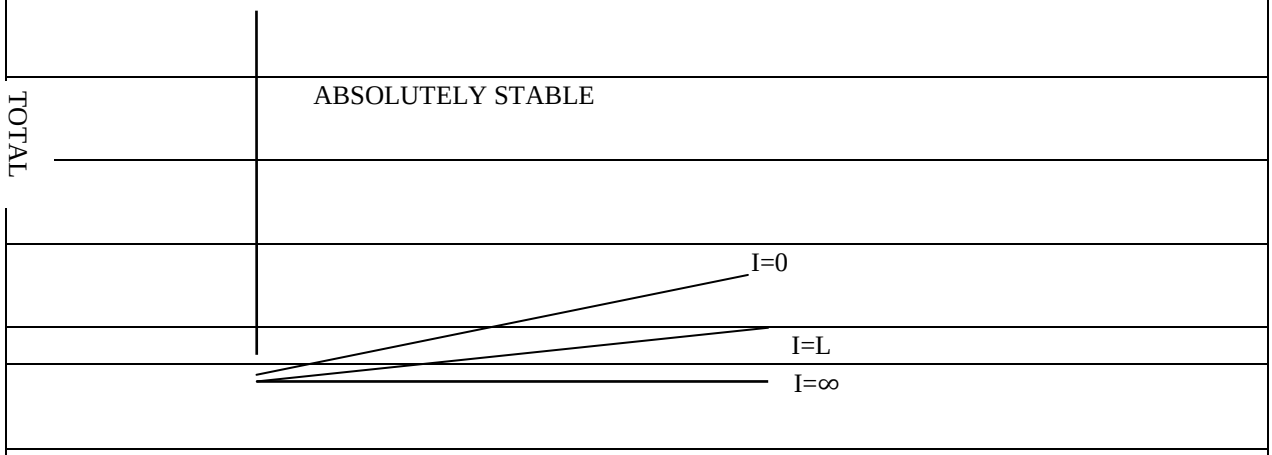
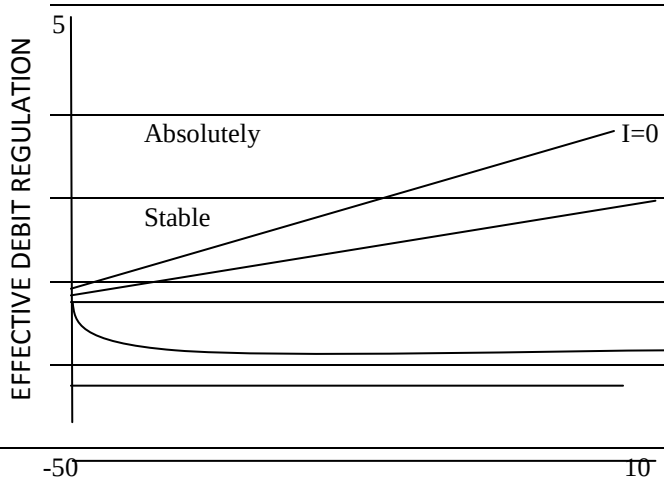



FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison?

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are

defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^\infty \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$ then given the constraints of $f(T_8)$, $(A1.5)$ *can never be satisfied Co*nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8 \lambda_8 e^{-T_8 \lambda_8} + [(U_8 - V_8) M_8 + Y_8] - [((U_8 -$$

V 8M 8–]**e**–T 8λ 8=0

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$\boldsymbol{Z}_8^2 + M_8 Z_8 \sin(Z_8 T_8) - (A_8 M_8 + Y_8) + (A_8 M_8 + X_8)\cos(Z_8 T_8) = 0$$

$$(A_8 M_8) Z_8 - M_8 Z_8 \cos(Z_8 T_8) + (A_8 M_8 - X_8)\sin(Z_8 T_8) = 0$$

Where $A_8 = U_8 - V_8$Using (A2.3), substitute for $\boldsymbol{Z}_8^2$ and $Z_8$in (A2.2):

$W_8^2 (X_8^2 - 3 M_8 A_8 X_8 - M_8^2 (X_8 - Y_8) + 2 M_8^3 A_8 + 2 A_8^2 M_8^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_8 = 1 - \cos(Z_8 T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8 / 2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_8 \le 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2\ (\mathrm{X}_8 - Y_8)$$

$$Q_8(2, A_8) = \left(A_8 + 2\,\mathrm{M}_8\right)^2 (2A_8\,\mathrm{M}_8 - (\mathrm{X}_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2\left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8\left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8\left(2X_8 + Y_8\right)\right)$$

$$+2A_8\left(X_8 + Y_8\right)$$

$$=2\,M_8^2 W_8^2 (M_8 + \mathrm{A}_8) + M_8\, A_8^2 W_8 + \left(2A_8 M_8 - X_8 + Y_8\right)(A_8 W_8 + M_8 W_8^2$$

$$+ (2- W_8)((2\,M_8\,X_8 + M_8\,Y_8)\,W_8 + A_8\,(X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $\mathrm{W}_8$touches the $W$ –axis only at its minimum value in which case (A2.9) shows that further increases in *A* will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range $0 \le W_8 \le 2$.Defining

$Q_8\left(W_8, A_8\right) = a_8 W_8^2 + c_8$ *where* $a_8\, b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \square b_8 \square 0$$

And

$$b_8^2 \square 4a_8c_8$$

$$Case(i): Y_8 = 0$$

$Let\ Y_8 = 0$ And try the substitution A= $X_8 = /2M_8 in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8 /$ $2M_8$) is positive for *W<2. Since* $\mathrm{Q}_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\ X_8$ ,

$Let\ Y_8 =\ y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 +\ n_8 y_8)$ we have

$$W_8^2 y_8\left[X_8 n_8 \frac{(1 + n_8 y_8)}{(2 + n_8 y_8)} + M_8^2\left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8 X_8 \left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2}+2y_8\left(n_8-1\right)+\frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)}=0.\ (A2.12)\right.$$

Since $n_8$ and $A_8$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8=(2+n_8 y_8),$ *the condition*(A2.88?)*can be written as*

$[\frac{X_8(E_8-8)}{E_8^2}]^2[\frac{X_8^2}{M_8^2}-4y_8 E_8^2\frac{X_8}{M_8^2}+4y_8(n_8-2)E_8^3]$ ?0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8/W_8^2)$ can be expressed as

$2E_8[y_8E_8+(y_8 E_8(8+y_8)(E_8-n_8))^{8/2}]>\frac{X_8}{M_8^2}>2E_8[y_8E_8-y_8E_8(8+y_8)(E_8-n_8))^{8/2}]$

**ABSOLUTE STABILITY**

Hence if $X_8/M_8^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0?$W_8$ ?2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8-8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2-2y_8\frac{(E_8+n_8^2y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right)-4y_8(E_8-n_8)\ ?0$$

And has the solution

$$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8-8)}[y_8(E_8 + n_8^2\, y_8) + [y_8^2\,(E_8 + n_8^2\, y_8)^2 + 4\, y_8 (E_8 - 8)(E_8 - n_8)]^{8/2}]$$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2\, E_8 [y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + (M_8 + U_8)\lambda_8 - (M_8 + V_8)\lambda_8 e^{2T_8\lambda_8} + (M_8 U_8 + Y_8)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8\, e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + (M_8 + V_8) Z_8 \sin(Z_8 T_8) - (U_8 M_8 + Y_8) + (M_8(U_8 + V_8) - X_8)$$

$$\mathrm{X}\cos(Z_8 T_8) - 2M_8 V_8 \cos^2(Z_8 T_8) + M_8 V_8 = 0$$

$$(M_8 + U_8) Z_8 - (M_8 + V_8) Z_8 \cos(Z_8 T_8) + (M_8(U_8 + V_8) - X)$$

$$\mathrm{X}\sin((Z_8 T_8) - M_8 V_8 \sin(Z_8 T_8) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 = 1 - \cos Z_8 T_8$and $A_8 = U_8 - V_8$we obtain

$$4M_8^2 V_8^2 W_8^4 + 2M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8(M_8 + V_8) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2[M_8 - V_8)^2(2A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8(M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8(A_8^2 + 4(X_8 + Y_8))]$$

$$+ W_8[A_8^2 + 2X_8)(2A_8 M_8 - X_8) + 2A_8 Y_8(M_8 + V_8) - M_8 A_8^3] + A_8^2(X_8 M_8 - Y_8) = 0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8 > 0$$

$$and$$

$$C_8 > (X_8 - Y_8)/2M_8$$

$recalling\ that\ A_8 = U_8$ and defining $C_8 = U_8 + V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8(0, U_8 V_8 = A_8^2 (X_8 + Y_8)$$

$$Q_8\left(2, U_8 V_8\right) = \left(C_8 + 2M_8^2\right)^2 (2C_8 M_8 + Y_8 - X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$the existence of a D-position at this point is determined by letting cos $(Z_8 T_8) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all *T*, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**ABSOLUTE STABILITY IN PREDATOR- PREY MODELS**

**APPLIED TO ACCENTUATION ATTRITION MODELS**

We apply model to each and every module and a concatenated model was to be presented which we procrastinate due to spatial restrictions.

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - D(R) - NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

The question is political-discursive. It's what Stengers and Latour are getting at with their concepts of diplomacy and cosmopolitics. They grant, first, that all entities exist and, second, that to say that someone's cherished idol (or whatever disputed entity they hold dear) is non-existent is a **'declaration of war'** – 'this means war,' as Stengers often says. They thus shunt onto-political discourse off of the terrain of knowledge/belief in the sense of existence/non-existence. Their basic claim seems to be that 'respect for otherness,' i.e. political pluralism, can only come from granting the entities that others hold dear an ontology, even if you don't 'believe' in them. You

are thus permitted to say 'I do not follow that god, he has no hold over me' but you are not permitted to say 'your god is an inane, infantile, non-existent fantasy, grow up.' And it's not just a question of politeness (although there's that too). The point is to grant others' idols and deities an existence – one needn't agree over what that existence entails, over what capacities that entity has or what obligations it impresses upon you as someone in its partial presence but to deny it existence entirely is to 'declare war' – to deny the possibility of civil discourse, of pluralistic co-existence. I believe that it was Richard Rorty who once quipped something like claims to reconcile realism and idealism always seem to end with a triumph of idealism. We don't, in fact, get realism through such approaches, but rather just get a pervasive anti-realism. I think this is also the problem with the "non-controversial pluralism" advocated by Stengers and Latour that Phillip defends here. Such pluralism is not realism but is, in fact, a thoroughgoing social constructivism. I think this is the central problem with Latour's argument in Irreductions (these days I regret having ever defended it). In rejecting both Enlightenment critique and what he calls "reduction" he wants to say something like "The Pentacostal really is filled with (e&eb) the Holy Spirit", that for the 19th working scientist heat really is a fluid and phlogiston really is what allows things to burn, and that for the Greek lightning really is an expression of Zeus's anger. Latour tells us that we aren't to reduce or explain away the entities posited by another group's "ontology" but are to develop explanations from within that ontology. It might not sound particularly sexy– and it certainly doesn't tell us what is worth thinking –but I can't help but believe that philosophy is the critical and reflective investigation of basic concepts that guide our investigation of the world about us, how we ought to live our lives, and what form of governance might be best. Compare two figures. A scientist might ask, what causes depression? We can very well imagine a philosopher turning around and asking the scientist, what is causality? The scientist presupposes a concept of causality in her investigations. She uses this concept in her inquiry. Now she might have a sophisticated concept of causality or she might never have thought much about causality at all, using it in the sort of colloquial and unreflective way that Plato decried when, for example, people

like Euthyphro talked about piety. A whole cascade of questions arise when we raise a question like "what is causality?" We can ask whether or not causality exists at all. We can ask how we distinguish between correlation and two events that merely accompany one another from genuine causation. This, for example, was Hume's question. But perhaps most importantly we can ask whether there is only one form of causality or many forms of causality. Is there only one-to-one causation; one cause and one effect? Is there many-to-one causation; or many events conspiring to produce an effect? Is there one-to-many causation; or one event producing a variety of different effects? We can even ask whether causality necessarily moves from past to present or whether there aren't forms of causality that move from future to past! **January 23, 2014 Pluralism and Realism:** multiple-worlds Over at Struggles Forever, Jeremy Trombley has an interesting post up on "the ontological turn" in anthropology or ethnography. I've been meaning to have a discussion with him about this as I think it's an issue many of us are struggling with. For example, the core project of The Democracy of Objects– a project which I think many have missed –is to somehow reconcile some version of social constructivism with a realist ontology capable of making room for ecology (which requires realist and materialist positions as there's a fact of the matter where global warming is concerned) as well as the role played by objective agencies in social assemblages such as technologies, infrastructure, features of geography, local climates, the growth cycles of plants and animals, waste, etc. Maybe we can try to organize some cross-blog event to discuss these issues. I certainly think they're close to the heart of Jeremy, Michael of Archive Fire, Arran James, and a host of others. As an aside, I'm beginning to realize how the different sites of the political I've been outlining– semiopolitics, thermopolitics, oikopolitics (political economy), geopolitics, eropolitics (the politics of sex and desire), biopolitics, and chronopolitics (and I'm sure there are other political sites!) –are drawing me away from traditional Marxism. Assuming that classical Marxism holds that economics or the conditions and relations of production are determinative of all other sites of the political, the various sites of the political that I've been outlining would lead to the conclusion that there is not one determinative

base of the political. This would not require committing Marx to flames, but rather of recognizing the phenomenon of overdetermination, or of a variety of different entangled sites of the political. But I digress. First, I find myself wondering what the ontological turn means in ethnography. Is it 1) the investigation of the different ontologies held/proposed by different cultures? E.g., the Aztecs believed that reality was structured in this way, while the Greeks in that way, and the ancient Chinese this way, etc? Or 2) is it an investigation of how real entities– independent of cultural beliefs –influence cultural formations? Or is it a combination of both? A position that I would favor. January 20, 2014 Model is applied to structured reality

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

**The evaporation spectrum of black holes from a local quantum gravity perspective Aurélien Barrau (Submitted on 21 May 2016**)

(7) Authors revisit the hypothesis of (e) a possible line structure in the Hawking evaporation spectrum of black holes, due to (e) non-perturbative quantum gravity effects, even arbitrarily far away from (e) the Planck mass.

(8) Authors show that this naive prediction might in fact hold in (eb)the specific context of loop quantum gravity, with (e&eb) a small departure from (e) the ideal case for (e) some low-spin transitions.

They also show that the effect is (=) neither washed out by the dynamics of (e) the process, nor by (e) existence of a mass spectrum up to (e) a given width, nor by the secondary component induced by (e) the decay of neutral pions emitted during (e&eb) the time-integrated evaporation. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Cite as: arXiv:1605.06613 [gr-qc] (or arXiv:1605.06613v1 [gr-qc] for this version

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t - T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\left(R(t)\right)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

It's very easy to treat this as an abstract, academic debate: "Are you a realist or are you an anti-realist?", as if it were just a matter of what happens to be true. But it seems to me that this debate has, in reality, always been about politics. As I outlined in my last post, we have perpetually seen how appeals to the real and natural have been used in the name of oppressive power, inscribing both the exploitation of nature and the oppression of various people in the very fabric of being itself. Theistic theology and realist ontology have perpetually been used in the name of

what Deleuze called "State Philosophies" or philosophies that ontologize contingent orders of power and privilege (e.g., "the great chain of being" used to justify patriarchy, monarchy, serfdom, poverty, etc, and appeals to nature used to justify poverty and racial inequality (The Bell Curve), patriarchy (evolutionary psychology), heteronormativity, capitalism, etc). Because arrangements of power and inequality are always contingent in the sense that there's no marked difference in the capacities of peoples, power always looks for a transcendent supplement that would provide justification through ontological necessity. Antirealism– from the Greek atomists to present –became the radical and emancipatory gesture because it revealed the lie behind all of these forms of social organization or their inherit contingency or arbitrariness. Realism, by contrast, has all too often functioned as an apologetics for arbitrary power and social organizations. Here it's worth recalling what Foucault said about science: "…Even before we know to what extent something like Marxism or psychoanalysis is analogous to a scientific practice in its day-to-day operations, in its rules of construction, in the concepts it uses, we should be asking the question, asking ourselves about the aspiration to power that is inherent in the claim to being a science. The question or questions that have to be asked are: "What types of knowledge are you trying to disqualify when you say that you are a science? What speaking subject, what discursive subject, what subject of experience and knowledge are you trying to minorize when you say 'I speak this discourse, I am speaking a scientific discourse, and I am a scientist.' What theoretico-political vanguard are you trying to put on the throne in order to detach it from the massive, circulating, and discontinuous forms that knowledge can take?" (Society Must Be Defended, 10). All of these questions hold equally for claims to something being real. What is one trying to minorize when claiming something is real? What becomes privileged? What is excluded? It is these questions that have been at the heart of the realism debates, for as Spencer-Brown taught us, every distinction has a marked and unmarked space, draws attention to something to be included and pushes something into the unconscious or the domain of the invisible, hidden, or veiled. This is above all the case with evocations of the real. However, as I've tried to show antirealism

leads to its own problems. First, so long as we exclude real beings from our ontological inventory, we are unable to fully understand how power functions ("such and such a set of cultural formations have been thoroughly debunked, yet people still live as if they believed them"). Not only do we not fully understand the sources of the problems due to too much focus on the discursive and semiotic, but we deny ourselves valuable sites of political intervention at the level of infrastructure. Second, we are prevented from addressing things such as cultural racism, such as that found in Heidegger's privileging of the West and the Greeks and Germans in particular. We do a good job addressing biological and theological racism, heteronormativity, and sexism by showing how it is a cultural construction, but when faced with racism such as Heidegger's where he argues that there's something "unique" about the Greek event, it's language, and about the German language, or Badiou/Zizek's racism with respect to the "Pauline Event", we really have no response. Here someone like Jared Diamond or Fernand Braud el is needed to explain global-geographical inequalities. Third, the tools of the cultural turn really do not provide us with the means of thinking the ecological as a site of the political. For this reason, I've tried to formulate a third way– which might be called "constructivist realism" or "constructivist naturalism" –that retains the insights of the cultural turn, while also allowing a robust place for the material. I don't claim to be original in this. I think that many such as Manuel DeLanda, Deleuze and Guattari, Stacy Alaimo, Karen Barad, Jane Bennett, etc., are up to something similar. It's a vast project that requires the work of a multitude of voices, especially given the way in which the culturalism pervades humanities. **The Relational and the Non-Relational: Notes towards an Immanent and Pluralist Theory of Meaning Posted by larval subjects under uncategorized k-bigpic I have deleted or reformulated some sentences for which I beg pardon attributable to spatial constraints** What is correlation, what is causation? What is coordination? What is coincidence? These thoughts have bothered humans since long. I have been able to excavate some reliable material from Wikipedia which I would like to share. "All views are partial and contingent – that's the lesson of pluralism, and you say as much yourself when you say that we must be attentive to ways our own

knowledge might contain superstitions. But it's not discourse that determines whether one view works and another doesn't – it's a confrontation with **non-discursive (non-human)** agencies." I wonder if part of the issue here is that we understand pluralism differently. For me this doesn't sound like pluralism at all. Rather, it just sounds like our **epistemological condition**. We [hopefully] come to understand those portions of existence we question or investigate. The reason we come to investigate them is largely contingent. Finally, the accounts of these features of existence we give can be mistaken. All of this is perfectly consistent with a monism. I'm not sure why what he outlines above is the lesson of pluralism. It seems to me that every realist knows this. It seems to me that ontological pluralism is something quite different. Ontological pluralism is the thesis that there are many different worlds inhabited by many different entities. Thus, for example, you would have one world; say that of Lucretius, that's only inhabited by atoms and their combinations. You would have another world; say the world of the Mongolian shaman, that's inhabited by spirits that do all sorts of things. The ontological pluralist is saying that these spirits are, that they exist, that they're real. This is something quite different than merely saying that there are partial and contingent points of view on the world. It's this saying that these entities are rather than that some person or groups of people believe that they are that is the nub of the issue. Now, I think part of the issue here is that there's an ambiguity in the term ontology. Ontology can be one of two things. On the one hand, ontology is a group or persons set of beliefs as to what is. Here it's trivially true that there are pluralities of ontologies and the realist readily recognizes this. This is the whole reason there are debates over ontology. Mongolian shamans have their ontology, Europeans theirs, Christian fundamentalists theirs, materialists theirs, etc. When striving to understand and communicate with others it's vital to understand these ontologies because, as rhetoricians like Burke point out, our **beliefs about what is** are among the things that motivate our action. Despite having never seen bacteria I was my hands because I believe there are bacteria and viruses on door handles and whatnot and don't want to get sick. My belief about a particular thing existing is what motivates my action. And who knows, perhaps this belief is as

superstitious as the belief that the crops failed because God was displeased with my community. On the other hand, ontology is a theory **about what is.** It is making a claim that something exists. This is where the rubber hits the road. The ontological pluralist seems committed to the thesis that every group set of beliefs about what exists is sufficient for granting the existence of those entities. This is what I find objectionable in Latour. Obviously I'm not bothered by Latour's suggestion that we should take into account the role that nonhumans like speed bumps, rivers, microbes, etc., play in the form that social assemblages take. But this is not all that Latour claims. This is precisely where the philosopher might balk. Unlike the ethnographer, the philosopher is not interested in what people believe exists, but rather philosophers– at least of the realist variant –are trying to figure out what is. In other words, the realist philosopher begins with the premise that not all of these beliefs about what is are true. So for the philosopher, recognizing that Mongolian shaman's believe in the existence of shamans would only be the first step. The next step would consist in determining whether there's good reason for thinking such entities really do exist, i.e., whether there's good reason for believing these entities have mind and culture independent reality. Lest readers think that I'm just picking on the supernatural here, we can ask similar questions about strings, subatomic particles, galaxies, etc., etc., etc. I guess the question really comes down to what exactly we mean by ontological pluralism. When we talk about ontological pluralism are we defending the thesis that people have different theories of being? If so, then ontological pluralism is trivially true. If this is what is meant, then I certainly share Jeremy's view that it's valuable to understand the different world's people believe in. Certainly when I was practicing as an analyst I didn't get in ontological debates with my patients and it was necessary to understand their theory of being or their ontology to properly attend to them. Or, when we talk about ontological pluralism are we defending the thesis that all these ontological theories are true and refer to really existing entities? That's quite a different claim and is not one I would defend or endorse. Now someone might object that "in both The Democracy of Objects and Onto-Cartography you defend the thesis that there are multiple worlds." This is true. Because I hold that not everything

is related I'm led to the conclusion that there are diverse worlds. However, I also hold that however many worlds there might be, these worlds are nonetheless composed solely of material entities. Within the framework I propose I wouldn't suggest that there's one world where there are spirits and another world where there are souls and yet another where there are only material entities. My view is that there aren't spirits or souls in any of these worlds. One of the things I keep hearing in these discussions is that somehow the realist adopts a view from nowhere. I honestly don't understand this criticism. Investigation always occurs somewhere and requires all sorts of mediations involving technologies, experiments, etc. It's that labor of gathering evidence, conducting experiments, using technologies to observe the world, etc., that gradually gives us a body of data that allows us to say there's good reason to believe that such and such a thing exists and has these powers. Another charge seems to be that the realist refuses to recognize that their claims about the world are fallible. I find this charge particularly strange because it's precisely because the realist recognizes the difference between our theories of the world or what we say about the world and the world itself that fallibility is built into the core of his position. Realism doesn't mean one holds they have special access to the world, that they know all truths, or that they have the truth in hand, only that there are truths to be known and that we can be mistaken about things. Circling **Squares Philip** has a post responding to my quandaries about how to mesh realism and pluralism. He writes: Ontologically and metaphysically the idea of realist pluralism is no longer an issue. There are (appropriately) numerous variants but the basic idea that reality is itself pluralistic is well established. Captioned expatiation is about Everett worlds and string theory

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from

being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

The von Neumann formulation leads us to conclude that the lower realms of quantum reality form a virtually limitless array of potentials that are inherently incapable of realization without the observation of a knowing entity, which we identify with consciousness (Kafatos and Nadeau, 1990). We thus have a correspondence between quantum theory, the hierarchical structure implied by the Holographic Principle, and the higher orders of experience,

leading to the full expression of the Conscious Universe. David Bohm's theory of the implicate and explicate orders involves a holographic Principle that is fully consistent with the Holographic Principle discussed here (Bohm 1980; 1986; Bohm and Hiley, 1993). According to Bohm, there is an implicate order that represents the universal, holographic subtext of reality, and which unfolds in every moment to produce the explicate order that we all observe. Thomas Germinario (2004) has equated the implicate order with (eb) the unconscious process, and the explicate order with (eb) conscious process. He emphasizes the nature of dreams within (eb) the implicate order, and the importance of the dream work for maintaining a healthy mind through (e&eb) integration of the implicate order and our daily lives. Allan Combs and Mark Holland (1990) connected the implicate order or holomovement (Bohm, 1980) with (e&eb) Carl Jung is theory of synchronicity (Jung and Pauli, 1955), with (e&eb) the implicate order providing (eb) a holographic medium through which apparently disconnected individuals become connected. The principle of synchronicity, the instantaneous connection of people and events beyond (e) the senses, has been equated with the quantum-physical principle of non-locality (Combs and Holland, 1990; Germine, 1991), and has been proposed to be the fundamental mechanism of conscious process (Germine, 1991).Applying the Holographic Principle theory of mind, subjects may be connected synchronously in the manner of two individuals conversing through cellular phones, where the radio signals are transmitted through (e&eb) a distant satellite. The difference here is that the radio signals travel at the speed of light, and thus the there is some miniscule time lapse between the sending and receiving ends. The Holographic Principle implies (eb) a much more distant and instantaneous communication pathway, with (e&eb) the analogue of the satellite being the holographic boundary of the Universe, as well as a much richer communication, potentially involving a transmission of thoughts and feelings. Our ego-consciousness seems to mask (e) the universal relatedness implied by (e) the Holographic Principle, and it is perhaps only through transcendence of (e) the ego-consciousness that the higher orders of experience can become conscious. In the phenomenon of synchronicity, there seems to be a meaningful connection between (e&eb) individuals that

breaks through the barrier of ego-consciousness. Such a connection is reported by many individuals in the course of dreams, when (e) the ego-consciousness has been suspended, at times informing (e&eb) the dreamer of something that has happened in the life of a meaningfully-connected individual. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Impassioned-intense Connectionist neuroscience, in an effort to attain relevance to clinical practice, has recently given us a theory of the relationship of the unconscious to (e&eb) consciousness which is (=) dualistic (Viamontes and Beitman, 2007), an ad hoc theory of conclusions based on (e) assumptions, the most erroneous of which are that representations of (e) information are information, and that such representations elaborate themselves through (e&eb) a process of representations of representations in neuronal circuits (Bechtel and Abrahamsen,1991). As theory, the connectionist models are lacking (e) heuristic value and **are (=) inimical to the progress** of (e) basic and clinical neuroscience. This is not to denigrate or minimize the importance of neuronal connections, but these connections are rather like the wires in (eb) a radio that plays a symphony. The radio, in and of itself, is (=) incapable of playing (e&eb) the symphony. The radio doesn't do not write the symphony, nor does it broadcast the Passionate-raging symphony over the airways. Yet, if we loosen one wire, the symphony is (=) not manifested through the radio. The brain is (=) not subject to (e) the same vulnerability, in that it has (e) parallel and redundant networks of processing, and, it is in this sense that the connectivity of the brain becomes (eb) important. Connectionism, to the extent that it accepts (e) consciousness and the (e&eb) unconscious at all, attributes them to (e) separate groups of "circuits" in the brain, with (e&eb) consciousness-processing circuits having (e) a limited but detailed ability to analyze (e&eb) information, and unconscious-processing processing more information in less detail (Viamontes and Beitman, 2007). There is an element of truth in this view in that information that has importance and meaning resonates within the brain at higher levels of experience and at higher levels of holographic recursion. This resonance involves what we know, what we feel, and what we are capable and willing to consciously realize. The Universe, as William Blake noted,

can be seen in a grain of sand. Indeed, the grain of sand requires (e) the Universe to exist. The human organism requires (e) the Self to exist. The Self is not a simple amalgam of (e) sensory experiences, as suggested in the connectionist model (Blinder, 2007). This amalgam view of the Self can be damaging when applied to (e&eb) psychotherapy, as the connectionists suggest (Blinder, 2007). The Self is ours forever and ever, from eternity to eternity. As Milton suggested, the mind is (=) its own place. Our deepest purpose is for our minds to resonate with (e&eb) the supra-conscious levels of experience, much as it is (=) purpose of the radio to play the symphony. If the symphony is not to our liking, (eb) we can turn the radio off. So, perhaps, the transpersonal consciousness can turn (e&eb) our own individual and collective radios off, or perhaps turn the volume down, if (e) we are not in harmony with the **Universal Mind**. Physics has given us many enigmas. Quantum theory itself is an enigma. But some of the enigmas seem to involve (e&eb) the limitation of (e) our own abilities of discernment. The thought experiment of (e) "Schrödinger's Cat," as it is often interpreted (e.g. Gribbin, 1984), assumes (eb) that the **cat is not conscious, and is (=) incapable of observation in the context of (e) collapse of (e) the wave function.** This is due, in large part, to the false dichotomy of (e) information and (e&eb) experience. Our observation of other animals indicates, to those that are attuned to (e) the animal mind, who can "feel" the animal mind, that, at (eb) the very least, all mammals are (=) conscious. It would seem that the great neuropsychiatrist, Stanley Cobb (1948), may have been correct in his attribution of (e) levels of consciousness to a wide variety of organisms: "...lower animals with (e&eb) no cerebrum appear to be conscious...even plants such as the sunflower that turns towards strong light may have (e) a vague awareness or warmth and comfort. There are many degrees of consciousness and it is my contention that it is integrated at (e&eb) many levels like other important functions of (e) the central nervous system."With respect to the theory of mind, recent mainstream thinking remains (eb) classical and mechanistic. Such views of mind are advocated by our most in prominent neuroscientists (e.g. Changeux, 1985; Gazzaniga, 1985) as neuroscience research attempts to explain all mental function on the basis of (e) **pure mechanism and the localized function of**

**specific areas of the brain.** Francis Crick, the distinguished Nobel Laureate, had a very influential second career as a neuroscientist. He concluded that we are basically soulless creatures, that all mental processes could be reduced to (e&eb) identifiable, neural correlates in (eb) the brain, and that, in fact, we are (=) nothing but a collection of neurons (Crick, 1995). Crick even went so far as to say that what we call the soul is (=) a group of neurons located in the prefrontal lobes of the brain. Another Nobel Laureate, Gerald Edelman (1987), has written a critically acclaimed book on "neural Darwinism," which purports that the brain is "circuitry" is constructed through a developmental process of "survival of the fittest" neurons and neuronal connections. At the same time the total purposelessness and lack of direction in evolution is expounded, again to critical acclaim, by Richard Dawkins (1986). The idea that there essentially no self or soul imbued with agency and self determination was expounded by Gilbert Ryle (1949/2002) with his parody of the "ghost in the machine," and this work continues to be influential among many scholars of the mind. Philosophers such as Patricia Churchland (1986), who coined the term "neurophilosophy," take an "eliminative materialist" view of (e) belief, free will, and consciousness. Other philosophers, such as Daniel Dennett (1991) would make (eb) experience "epiphenomenal," a by-product of (e) brain processes with (e&eb) no effect whatsoever on (eb) the function of mind. However, unlike neuroscience, which is still in its infancy, the philosophy of mind has had a long history, including such notable figures as Plato, Aristotle, Descartes, Kant, Locke, Hume, Kierkegaard, and James, to name just a few. Perhaps, as noted by Ervin Laszlo (1974): "in today's world, most of the traditional functions of (e) cognitive synthesis have atrophied (e) and are ignored and neglected." Hopefully, what we have presented here will be a step in the right direction for what Laszlo calls a "conceptual synthesis" to help fill the need for (e) meaningful engagement is such a world. The fundamental elements of such a synthesis have been described by (e) Laszlo as follows: Conceptual synthesis performs (e&eb) at least five basic functions in the guidance of affairs. They are the mystical, the cosmological, the sociological, the pedagogical or psychological, and the editorial functions. The mystical function inspires (eb)

in man a sense of mystery and profound meaning related to (e&eb) the universe and of himself in it. The cosmological function forms (eb) images of the universe in accord with (e&eb) local knowledge and experience, enabling (eb) men to describe and identify (e&eb) the structure of (e) the universe and the forces of nature. The sociological function validates supports and enforces (eb) social order, representing (eb) it in accord with (e&eb) the nature of the universe, or as the natural or right form of (e) social organization. The pedagogical or psychological function guides (eb) individuals through stages of life, teaching ways of understanding themselves and others and presenting (eb) desirable responses to life is challenges and trials. Finally, the editorial function of (e) conceptual is to define (eb)some aspects of reality as important and credible and hence to be attended to, and other aspects unworthy of serious attention. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Kant sums up this inversion and its spirit early in the Critique of Pure Reason: Up to now it has been assumed that all our cognition must conform to the objects; but all attempts to find something about them a priori through concepts that would extend our cognition have, on this presupposition, come to nothing. Hence let us once try whether we do not get farther with the problems of metaphysics by assuming that the objects must conform to our cognition, which would agree better with the requested possibility of an a priori cognition of them, which is to establish something about objects before they are given to us. This would be just like the first thoughts of Copernicus, who, when he did not make good progress in the explanation of the celestial motions if he assumed that the entire celestial host revolves around the observer, tried to see if he might not have greater success if he made the observer revolve and left the stars at rest. (B xvi) Onticology, like all variations of object-oriented ontology, is realist in its orientation. In defending a realist ontology onticology holds that the vast majorities of objects, actants, beings, or entities are independent of humans and are what they are regardless of whether any humans regard them or register them. In short, onticology rejects any anthropomorphic, idealist, or anti-realist thesis to the effect that to be is to be the correlate of mind, spirit, the body, the human, and language or otherwise. While it is certainly the case that knowledge is necessarily dependent on the object to which it relates, the converse does not hold true. Objects are not dependent on being known, regarded, perceived, or spoken

about. As such, and to put it in Aristotlean terms, knowledge is an accident of objects, not objects an accident of knowledge. As Althusser so nicely puts it, "[n]o doubt there is a relation between thought-about-the-real and thisreal, but it is a relation of knowledge, a relation of adequacy or inadequacy of knowledge, not a real relation, meaning by this a relation inscribed in that real of which the thought is the (adequate or inadequate) knowledge" (Reading Capital, 96). Althusser goes on to remark that "[t]he distinction between a relation of knowledge and a relation of the real is a fundamental one: if we did not respect it we should fall irreversibly into either speculative or empiricist idealism" (ibid.). Onticology categorically endorses Althusser's verdict. It is a fundamental necessity to distinguish between those relations that belong to the object and those that belong to knowledge. Contemporary philosophy continuously confuses these two very different sorts of relations. Naturally the question arises of how it is possible to surmount our relation to the object so as to determine whether objects themselves possess the properties we encounter in relating to objects. In other words, given that we can only ever relate to the object in relating to the object how is it possible to surmount this relation to get at the being of the object itself? Much more will have to be said about this later– and the answers will be surprising with respect to standard prejudices about realism –however, for the moment it can be said that onticology takes its epistemological inspiration from the transcendental realism of Roy Bhaskar. Among other things, Bhaskar sought to provide a transcendental grounding for the sciences. Insofar as onticology defends the thesis that the field of being is much more vast than the field of objects investigated by the natural sciences, it parts way with the thesis that the domain of being is exhausted by the domain of natural objects. However, the general form of Bhaskar's argument holds for our realist purposes. A transcendental argument seeks to elucidate the conditions under which certain acknowledged practices and forms of cognition are possible. Kant, for example, asked what must be the case for mathematical judgments to be possible. How is it both that we are able to extend our knowledge, as if by magic, through mathematical judgments and, more significantly, that these judgments are able to provide genuine knowledge of the world despite the fact that these

forms of reasoning are not based on experience? Part of Kant's argument consisted in claiming that mind imposes the forms of space and time on the data of experience. In other words, space and time are not attributes of being itself but rather of the mind that regards being. Insofar, Kant argues, as mathematics is ultimately a rumination on the nature of space and time taken in their most abstract form and insofar as the mind imposes space and time on the manifold of sensation, it thus follows that a priori judgments about the nature of spatio-temporal relations are possible that anticipate the structure of actual-space times without directly experiencing these space-times. Why? Because any manifold of sensation must necessarily be structured by these forms imposed by intuition. **Onticology–A Manifesto for Object-Oriented Ontology Part I Posted by larval subjects under Object-Oriented Philosophy My Question Is This: You Say, "Granted There Are Deeper And Deeper And More Primitive Forms Of "Knowledge" Or Abstraction Into Mechanism, But I Wouldn't Say That The Atomic Level Brings A New Epistemic World Into View." But Doesn't Any Kind Of Knowledge, However Primitive, Imply Bringing Some Kind Of Epistemic World Into View Or Proto-View?** Joel: David, Yes Exactly. My Question Is, Just How Far Down The Complexity Gradient Can The Meaning Of 'Knowledge,' Or 'Semiotics, Or Specifically Representation (Maya) Be Stretched Without Breaking? I Can Comfortably Stretch It To The Genetic Level, As The Proto-Epistemic, Where Genetic Codes 'Represent' Phenotypic Structures Or Modes, And The Whole Of Evolution Learns From Its Mistakes And Successes As That Knowledge And History Is Encoded Into And As The Recapitulation Of The Embryogenesis Of The Organism. But An Atom, In My Model, While Certainly Reacting To Its Environment From Its Own Attractor, Essence, And 'Interiority' (Prehension) Isn't Making Representations Of It To Do So. But The Key To The Power And 'Essence' Of Representation (And Simultaneously Its Limitation) Is The Capacity Of Abstraction And Choice. Evolution Does This At A Rudimentary Level (And Anthropomorphizing A Bit) With Its Sending Into The World Its 'Random' Variations, Organismic Ideas, In A Sense. Its 'Mind' Is Literally On The Outside As Our Living World And Biosphere. A 'Choice' Is Made In The Biospheric Interiority,

And Knowledge For The Evolutionary/Embryogenetic Trajectory Is Gained By The "Differential Reproductive Success" Of These Individuals Lives Through Time, As The Increasing Intelligence, Representation, And Capacity For Choice Is Directly Injected Into The Flow Through Sexual Selection. But Where Is The Encoded/Abstracted (Enfolded) Representation Of Options And A Choice Among Them Made At The Atomic Level? I Am Working With The Most Cutting-Edge And Coherent Models For The Atom (Including My Own) That I Can Find Among The Heterodoxy, And I See No Evidence Of Representation And Choice At This Level, Or Rather No Real Way To Further Stretch The Meaning Of The Terms To This Level. Sure There Is Uncertainty From Infinite Difference And Immanent Causation (See Bohm On Infinite Causation), But Without Representation There Really Is No Choice (Although Crudely We Could Call A 'Bifurcation' A 'Choice' By The System As A Whole). So I Am Simply Saying That Representation At The Atomic Level Is Entirely Enfolded And These Are Really Just Ontic Level Phenomena. The Atomic Level, In My View, Is Pre- Not Yet Proto-Epistemic, Which I Place at the Genetic Level Where We Find the First Codes And (e&eb) Primitive Representation. This Just Means That There Is A Real Distinction Between Ontic And Epistemic (And A Gradient Between Them) And We Can't Collapse The Two Into Just The Ontic-Epistemic. And Yet, With The Distinction Intact, And The Gradient To Explore, They Remain Nondual, Univocal, ONE Or Ontic. What I Would Say Exists At All Levels (My Own Self-Aware Myth About The Given) Is What I Call The "Symbiogenesis Of Subject And Object". This Is The Spinozan Essence Of Dynamic Stability And Growth, And Leibnizian Prehension. Also Called The "Nucleation Of Observability" In Spinbitz. It Is From This Nucleus Of The Observer (Fuller) In A "Point-Free Geometry" (Whitehead) Of Pointless "Points Of View" As Conditions Of Boundary That Wilber's "All Is Perspective" Finds Its Ontic Grounding. And This Symbiogenesis Is Also A Rudimentary Or Deep-Level Native Or Primitive Intelligence, Evolution At The Involution Of The Subject-Object Interface. It Is An Exploration Of Creative Learning And Direct Immanent Awareness, Unmediated (Enfolded) By Representational Forms (Unfolded). It Just Is, In Waves (Cosmic Vertebrae)

Before It Abstracts, Represents And Incarnates The Flesh Of What Could Be. Matter Is Enfolded Seed Of The Flower Of Its Abstraction. Maya Is This Depth And Recursion Of Brahma Into And Through The Boundary Conditions Of Abstract Relation. And Brahma Is The Span Of The Emergent Rigor Of The Cosmic Ergodic Spine That Opens The Recursions Into The Field Of Infinite Difference. In Immanent-Transcendent Waves, Literally And Empirically A Limit-Cycle In Ergodicity, It Enfolds And Unfolds From Anatom To Anatomy, From Infinite Intelligence Or Omnirelational Awareness Of Self, Through The Representational Depths Of Self As Self-Consciousness. Omni-Evolution, Trans-Dynamics, And The Anatom Posted On February 21, 2013A Discussion With Tom Huston, Part II The Following Is Part 2 Of An Excerpt From A Discussion Initiated By Tom Huston On His Facebook Page, Edited And Embellished For Clarity. Find Part One, Here. In This Part We Dig Into The Nature Of The "Anatom" And The "Ergodic Spine" Of An Electro-Fractal Cosmos, And Into The Trans-Dynamic Integration Between Being And Becoming, And The X-Interface−The Crossroads Of The Ontic-Epistemic (Brahma And Maya) And The Subject-Object Polarities. Univocal_Dynamics_V1 Tom Huston Is A Founding Member Of Integral Institute And A Former Editor Of Enlighten Next Magazine. His Webpage Is Tomhuston.Com. David Marshall Is A Writer And Editor Living In Chicago. No meaning, value, or normativity to be found in nature, that there is nothing **natural beings ought to be,** but that, rather, these judgments arise from us (See last post of larval subjects) which has led some to raise valuable questions about the coherence of these claims. The problem arises when the following three propositions are taken together: 1. there is nothing outside of nature.2. Beings have no intrinsic meaning, purpose, or value (in and of themselves, there's nothing they ought to be). 3. Value judgments about what beings and being ought to be arise from us and beings like us (bonobo apes, dolphins, institutions, birds of paradise, etc). The problem arises between proposition 1 and 3. How can it both be true that there is nothing outside of being and those normative judgments belong to us and other beings capable of making normative judgments, not nature? The problem arises from restricting these judgments to humans and beings capable of

making these judgments. In making such a claim it seems as if we're saying that there's something outside of nature, something that is beyond nature, thereby violating the first thesis and potentially reintroducing the nature/culture distinction. By contrast, if we say that normative judgments are the special domain of those living beings with the proper degree of sentience to make such judgments, then we seem to reintroduce the nature/culture distinction and fall back into the sorts of problems that I outlined in my last post (and that are so nicely critiqued by thinkers such as Latour). Is there a way out of this? I don't know. Great_Chain_of_Being_2_ (lighter) ASIDE: It seems to me that this is really what the debate between realism and anti-realism, realism and socio-linguistic constructivism surrounding the new materialisms and speculative realists has really been all about.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \qquad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9}+\frac{N_9 dG_9}{dR_9}-\frac{dB_9}{dR_9}>\frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9}+\frac{dM_9}{dR_9}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usance bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_9=\frac{N_9 G_9}{2M_9}\frac{d(9nF_9)}{dR_9}(\geq 0) \qquad (9a)$$

$$I_9=-\frac{N_9 G_9}{2M_9}\frac{d(9nM_9}{dR_9}(\geq 0) \qquad (9b)$$

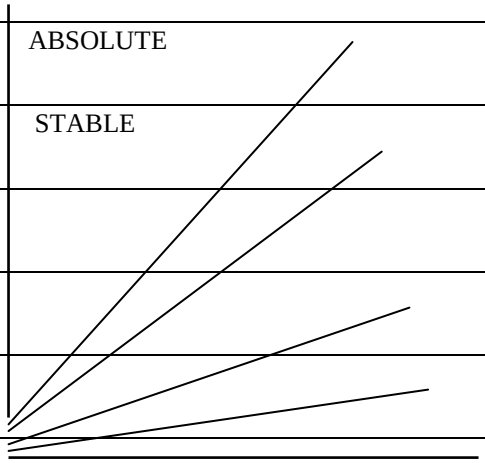
It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9>L_9-I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_9>L_9-I_9, \quad \text{when} \quad P_9^2/4>I \qquad (8a)$$

$$P_9>L_9-I_9+\frac{(I_9-P_9^2}{(L_9-I)} \quad \text{when} \quad P_9^2/4<I. \qquad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Reffer couffesy Nunney

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability

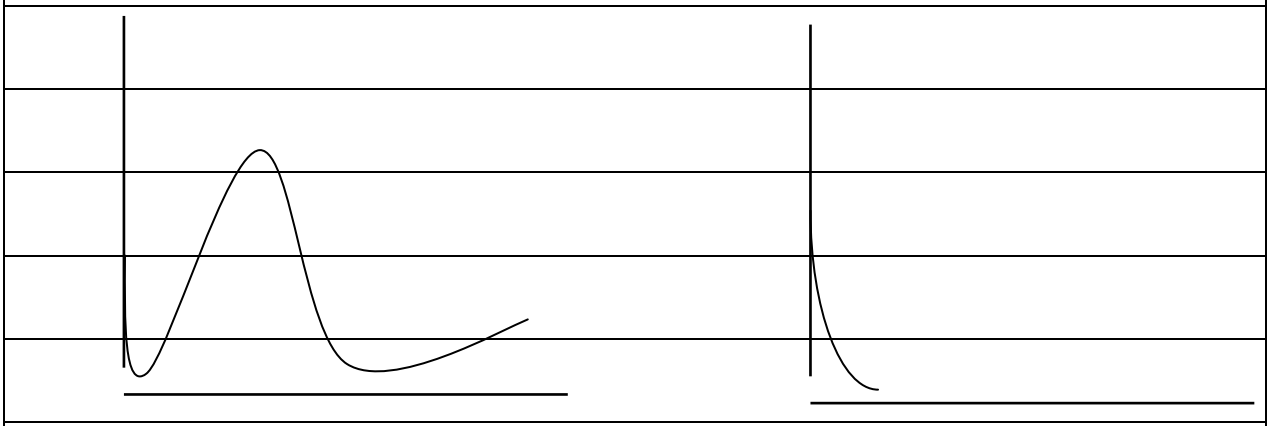
 

is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

Indeed, Nunny's model can be applied to some of the consciousness and philosophical problems. Predator prey analysis has ripped open the veneer of the fundamental concept of universe namely Debit-Credit world.

The noise of ordinary consciousness and (e&eb) beliefs drowns out (e) the signal. Unconsciousness is the background of (e) our ordinary awareness. Our organism is very much at (eb) the center of such effects. The organismic source is (=) our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is a meme-Scape. Pre-conceived concepts vie with structures, concepts with images. Like scientists who ignore assumed truths, we leapfrog over our beliefs and personality deficits, claiming idiosyncratic imagination is literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is both. We inherit the structures that make our experience what it is.

But the structure itself is “empty,” and each human culture “fills” it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us

feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through 're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns, nomenclature, and internal coherence of meaning and representation. In a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist

simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** The lower domain of experience is what we do not directly observe, the quantum realm, while the higher level is what we ordinarily observe, the classical realm. Based purely on experiment, the formulation of quantum theory initially placed the great divide between the quantum and classical domains between the measuring apparatus and the particle (Herbert, 1985). However, in the early 1930s, John von Neumann, in his rigorous mathematical treatise on quantum mechanics, found no support for such a division, compelling him to conclude that the wave for was collapsed by consciousness. Von Neumann's treatise has been called the "quantum bible" and "the most influential book on quantum theory ever written" (Herbert, 1985).

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t-T_9)F_9\big(R_9(t-T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t-T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ and the effective prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting $dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

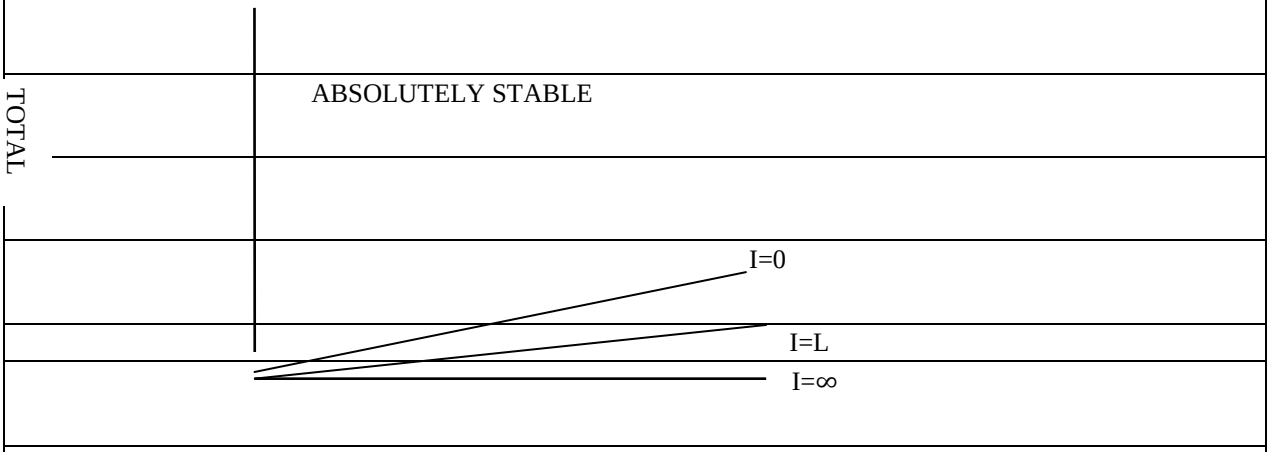



FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.

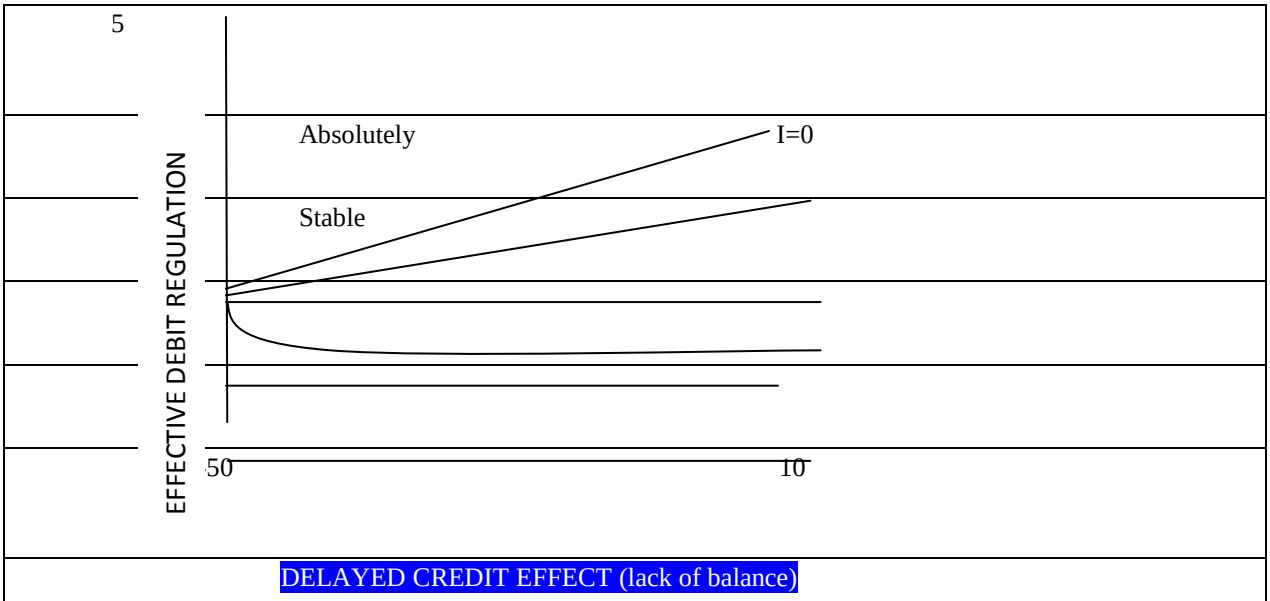



Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the

constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^\infty \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9\left(U_9 - V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9\right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$ then given the constraints of $f(T_9)$, $(A1.5)$ *can never be satisfied* Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDERATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 - V_9M_9 -]e^{-T_9\lambda_9} = 0$$

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9Z_9\sin(Z_9T_9) - (A_9M_9 + Y_9) + (A_9M_9 + X_9)\cos(Z_9T_9) = 0$$

$$(A_9M_9)Z_9 - M_9Z_9\cos(Z_9T_9) + (A_9M_9 - X_9)\sin(Z_9T_9) = 0$$

Where $A_9 = U_9 - V_9$ Using (A2.3), substitute for $Z_9^2$ and $Z_9$ in (A2.2):

$W_9^2(X_9^2 - 3M_9A_9X_9 - M_9^2(X_9 - Y_9) + 2M_9^3A_9 + 2A_9^2M_9^2]$

$$+W\left[MA^{3}-A^{2}\mathrm{X}+2\mathrm{MA}(2\mathrm{X}+\mathrm{Y})-2X^{2}\right]+A^{2}(\mathrm{X}+\mathrm{Y})=0$$

*where* $W_9$=1-cos ($\mathrm{Z}_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (\mathrm{X}_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ \mathrm{X}_9 - Y_9/2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range 0⍰$W_9$ ⍰2 *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (\mathrm{X}_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\ \mathrm{M}_9)^2\ (2A_9\mathrm{M}_9 - (\mathrm{X}_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9M_9^2 - 3M_9X_9\right) + W_9\left(3M_9A_9^2 - 2A_9X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9 + Y_9\right)$$

$=2M_9^2W_9^2(M_9 + A_9) + M_9A_9^2W_9 + \left(2A_9M_9 - X_9 + Y_9\right)(A_9W_9 + M_9W_9^2$

$+ (2- W_9)((2M_9X_9 + M_9Y_9)W_9 + A_9(X_9 + Y_9))$,

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range 0≤ $W_9$ ≤2.Defining

$Q_9(W_9, A_9) = a_9W_9^2 + c_9$ *where* $a_9 b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

**Harald Atmanspacher, Hans Primas** show that such an epistemic perspective can be consistently complemented by Einstein's ontically oriented position. **Time, Quantum and Information 2003, pp 301-321 Epistemic and Ontic Quantum Realities Harald Atmanspacher, Hans Primas Robert W. Spekkens** present a toy theory that is based on a simple principle: the number of questions about the physical state of a system that are answered must always be equal to the number that are unanswered in a state of maximal knowledge. Many quantum phenomena are found to have analogues within this toy theory. These include the noncommutativity of measurements, interference, the multiplicity of convex decompositions of a mixed state, the impossibility of discriminating nonorthogonal states, the impossibility of a universal state

inverter, the distinction between bipartite and tripartite entanglement, the monogamy of pure entanglement, no cloning, no broadcasting, remote steering, teleportation, entanglement swapping, dense coding, mutually unbiased bases, and many others. The diversity and quality of these analogies is taken as evidence for the view that quantum states are states of **incomplete knowledge** rather than states of reality. A consideration of the phenomena that the toy theory **fails to reproduce, notably, violations of Bell inequalities** and the existence of a Kochen-Specker theorem, provides clues for how to proceed with this research program.DOI: http://dx.doi.org/10.1103/PhysRevA.75.032110 **Evidence for the epistemic view of quantum states A toy theory Phys Rev A 75, 032110 – Published 19 March 2007 Robert W. Spekkens In such specialized and autonomous areas, it is to the eminent and erudite authors to interpret the contradistinctions and contradictions' in the dove tailing explaining. This I feel would be yeoman service in the creation of Nature's general ledger. M. S. Leifer and O. J. E. Maroney** examine the relationship between quantum contextuality (in both the standard Kochen-Specker sense and in the generalized sense proposed by Spekkens) and models of quantum theory in which the quantum state is maximally epistemic. We find that preparation noncontextual models must be maximally epistemic, and these in turn must be Kochen-Specker noncontextual. This implies that the Kochen-Specker theorem is sufficient to establish both the impossibility of maximally epistemic models and the impossibility of preparation noncontextual models. The implication from preparation noncontextual to maximally epistemic then also yields a proof of Bell's theorem from an Einstein-Podolsky-Rosen-like argument. DOI: http://dx.doi.org/10.1103/PhysRevLett.110.120401 Received 25 August 2012Published 20 March 2013 © 2013 American Physical Society **Maximally Epistemic Interpretations of the Quantum State and Contextuality Phys. Rev. Lett 110, 120401 – Published 20 March 2013 M. S. Leifer and O. J. E. Maroney.** Epistemological, onitic and ontological aspects particularly of quantum mechanics have invited

profound thinkers to contribute their mite to the growing literature on the subject, which includes the contributions from the subterranean realm and ceratoid dualism of philosophy like that Chaos theory of Deleuze, Derrida , Kierkegaard and others. **Andrei Khrennikov** shows that the Dirac-von Neumann formalism for quantum mechanics can be obtained as an approximation of classical statistical field theory. This approximation is based on the Taylor expansion (up to terms of the second order) of classical physical variables – maps f: Ω → R, where Ω is the infinite-dimensional Hilbert space. The space of classical statistical states consists of Gaussian measures ρ on Ω having zero mean value and dispersion σ2 (ρ) ≈ h. This viewpoint to the conventional quantum formalism gives the possibility to create generalized quantum formalisms based on expansions of classical physical variables in the Taylor series up to terms of nth order and considering statistical states ρ having dispersion σ2(ρ) = h n (for n = 2 we obtain the conventional quantum formalism). **December 2005, Volume 18, Issue 7, pp 637-650 Date: 28 Nov 2005 Generalizations of Quantum Mechanics Induced by Classical Statistical Field Theory Andrei Khrennikov** It is widely accepted that consciousness or, in other words, mental activity is in some way correlated to the behavior of the brain or, in other words, material brain activity. Since quantum theory is the most fundamental theory of matter that is currently available, it is a legitimate question to ask whether quantum theory can help us to understand consciousness. Several approaches answering this question affirmatively, proposed in recent decades, will be surveyed. It will be pointed out that they make different epistemological assumptions, refer to different neurophysiological levels of description, and adopt quantum theory in different ways. For each of the approaches discussed, these imply both problematic and promising features which will be indicated. **Discrete Dynamics in Nature and Society Volume 2004 (2004), Issue 1, Pages 51-73 http://dx.doi.org/10.1155/S102602260440106X Quantum theory and consciousness: an overview with selected examples Harald Atmanspacher Copyright © 2004 Hindawi Publishing**

**Corporation. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.** Another thought provoking and investigatory article is by Peter G. Lewis et al. Perhaps the quantum state represents information about reality, and not reality directly. Wave function collapse is then possibly no more mysterious than a Bayesian update of a probability distribution given new data. **Peter G. Lewis et al** consider models for quantum systems with measurement outcomes determined by an underlying physical state of the system but where several quantum states are consistent with a single underlying state—i.e., probability distributions for distinct quantum states overlap. Significantly, they demonstrate by example that additional assumptions are always necessary to rule out such a model. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.150404 **Distinct Quantum States Can Be Compatible with a Single State of Reality Phys. Rev. Lett. 109, 150404 – Published 9 October 2012 Peter G. Lewis et al** Communication complexity of a quantum channel is the minimal amount of classical communication required for classically simulating a process of state preparation, transmission through the channel and subsequent measurement. It establishes and perpetuates a limit on the power of quantum communication in terms of classical resources. **Alberto Montina** shows that classical simulations employing a finite amount of communication can be derived from a special class of hidden variable theories where quantum states represent statistical knowledge about the classical state and not an element of reality. This special class has attracted strong interest very recently. The communication cost of each derived simulation is given by the mutual information between the quantum state and the classical state of the parent hidden variable theory. Finally, find that the communication complexity for single qubits is smaller than 1.28 bits. The previous known upper bound was 1.85 bits. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.110501 **Epistemic View of Quantum States and**

**Communication Complexity of Quantum Channels Phys. Rev. Lett 109, 110501 – Published 12 September 2012 Alberto Montina for models see elsewhere .** Stability analysis of what? Quantum state representing reality or our knowledge of reality? Does the quantum state represent reality or our knowledge of reality? In making this distinction precise, **Nicholas Harrigan, Robert W. Spekkens** are led to a novel classification of hidden variable models of quantum theory. They show that representatives of each class can be found among existing constructions for two-dimensional Hilbert spaces. Approach also provides a fruitful new perspective on arguments for the nonlocality and incompleteness of quantum theory. Specifically, **Nicholas Harrigan, Robert W. Spekkens** show that for models wherein the quantum state has the status of something real, the failure of locality can be established through an argument considerably more straightforward than Bell's theorem. The historical significance of this result becomes evident when one recognizes that the same reasoning is present in Einstein's preferred argument for incompleteness, which dates back to 1935. This fact suggests that Einstein was seeking not just any completion of quantum theory, but one wherein quantum states are solely representative of our knowledge. Hypothesis is supported by an analysis of Einstein's attempts to clarify his views on quantum theory and the circumstance of his otherwise puzzling abandonment of an even simpler argument for incompleteness from 1927. **February 2010, Volume 40, Issue 2, pp 125-157 Date: 09 Jan 2010 Einstein, Incompleteness, and the Epistemic View of Quantum States Nicholas Harrigan, Robert W. Spekkens** statistical model of the probabilistic description of physical reality has been studied by many authors. **Andrei Khrennikov posits** a contextual statistical model of the probabilistic description of physical reality. Here contexts (complexes of physical conditions) are considered as basic elements of reality. There is discussed the relation with QM. He proposes a realistic analogue of Bohr's principle of complementarity. In the opposite of the Bohr's principle, **Andrei Khrennikov** principle has no direct relation with mutual exclusivity for

observables. To distinguish his principle from the Bohr's principle and to give better characterization, he changes the terminology and speaks about supplementarity, instead of complementarity. Supplementarity is based on the interference of probabilities. It has quantitative expression trough a coefficient which can be easily calculated from experimental statistical data. There is need not appeal to the Hilbert space formalism and noncommutativity of operators representing observables. Moreover, in our model there exists a pair of supplementary observables which cannot be represented in the complex Hilbert space. There are discussed applications of the principle of supplementarity outside quantum physics. **October 2005, Volume 35, Issue 10, pp 1655-1693 The Principle of Supplementarity: A Contextual Probabilistic Viewpoint to Complementarity, the Interference of Probabilities and Incompatibility of Variables in Quantum Mechanics Andrei Khrennikov Foundations of Physics** Description of the vacuum in Yang-Mills theory and arrive at a physical interpretation of the pseudoparticle solution and the attendant violation of symmetries has been done by **R. Jackiw and C. Rebbi.** It is here vaccuum stability forms the bastion, pillar, post and stylobate of the mechanisms of yang mills Theory and action of a particle. The existence of topologically inequivalent classical gauge fields gives rise to a family of quantum mechanical vacua, parametrized by a CP-nonconserving angle. The requirement of **vacuum stability** against gauge transformations renders the vacua chirally noninvariant. DOI: http://dx.doi.org/10.1103/PhysRevLett.37.172 **Vacuum Periodicity in a Yang-Mills Quantum Theory Phys. Rev. Lett 37, 172 – Published 19 July 1976 R. Jackiw and C. Rebbi.** Three-dimensional Yang-Mills and gravity theories augmented by gauge-invariant mass terms are analyzed. These topologically nontrivial additions profoundly alter the particle content of the models and lead to (eb) quantization of a dimensionless mass-coupling-constant ratio. The vector field excitations become (=) massive, with (e&eb) spin 1 (rather than massless with spin 0), and the mass provides (eb) an infrared cutoff. The gravitation

acquires (e) mass, mediates (e&eb) finite-range interactions, and has (e) spin 2 (rather than being absent altogether); although its mass term is of (e)third derivative order, there are no ghosts or acausalities. DOI: http://dx.doi.org/10.1103/PhysRevLett.48.975 © 1982 The American Physical Society **Three-Dimensional Massive Gauge Theories Phys. Rev. Lett 48, 975 – Published 12 April 1982 S. Deser et al** All scattering amplitudes in (eb) the maximally supersymmetric N=4 super-Yang–Mills theory possess (e) a new, dual superconformal symmetry which extends the previously discovered dual conformal symmetry of MHV amplitudes. To reveal this property authors formulate the scattering amplitudes as (=) functions on (e&eb) the appropriate dual superspace. Rewritten in this form, all tree-level MHV and next-to-MHV amplitudes exhibit manifest dual superconformal symmetry. Authors propose a new, compact and Lorentz covariant formula for (e) the tree-level NMHV amplitudes for (e) arbitrary numbers and types of external particles. The dual superconformal symmetry is broken at loop level by infrared divergences. However, authors provide evidence that the dual conformal anomaly of (e) the MHV and NMHV superamplitudes is (=) the same and, therefore, their ratio is (=) dual conformally invariant. Authors show this explicitly for (e) the six-particle amplitudes at one loop. Authors conjecture that these properties hold for all, MHV and non-MHV, superamplitudes in N=4 SYM both at (eb) weak and at strong coupling. **Nuclear Physics B Volume 828, Issues 1–2, 21 March 2010, Pages 317–374 Dual superconformal symmetry of scattering amplitudes in N=4 super-Yang–Mills theory J.M. Drummonda DOI: 10.1016/j.nuclphysb.2009.11.022.** By comparison with numerical results in (eb) the maximal Abelian projection of lattice Yang–Mills theory, it is argued that the nonperturbative dynamics of (e) Yang Mills theory can be described by (e) a set of fields that take their values in (eb) the coset space SU(2)/U(1). The Yang–Mills connection is parameterized in (eb) a special way to separate (e&eb) the dependence on the coset field. The coset field is then regarded as (=) a collective variable and a method to obtain its effective

action is developed. It is argued that the physical excitations of (e) the effective action may be (=) knot solitons. A procedure to calculate the mass scale of knot solitons is discussed for (e) lattice gauge theories in the maximal Abelian projection. The approach is extended to (e&eb) the SU (N) Yang–Mills theory. A relation between the large N limit and the monopole dominance is pointed out. **Physics Letters B Volume 458, Issues 2–3, 8 July 1999, Pages 322–330 An effective action for monopoles and knot solitons in Yang–Mills theory Sergei V. Shabanov DOI: 10.1016/S0370-2693(99)00612-7.** It is here the quintessential formulation of consciousness occurs. Like a man standing on the threshold of infinity trying to ponder what lies beyond the veil, which separates the seen from unseen a true seeker tries to lift the individual general ledger to the cosmic general ledger. Like in Darwinian evolution, mind also has to evolve from individual consciousness to cosmic consciousness. In Kaishmiri Shaivism one withdraws from the finite to the infinite, but one also goes on an outward journey from the infinite to the finite, because both the finite and the infinite have an intimate connection. The finite is not seen as unreal, but as a symbol of the infinite. There is no real distinction between them. Those two movements constitute Spanda, a key concept in Kaishmiri Shaivism. Spanda is the pulsation of the Absolute in different phases of being. There are no opposites like subject and object, unity or duality, absolute or relative. They are just different phases of the universal vibration of the Absolute. The goal is to realize or be (=) at once infinite and finite. One does not turn away from appearances (like in Advaita Vedanta), but one realizes that the Absolute manifests all things. Spanda, the eternal pulsation of the Absolute, oscillates between a passion to create and dispassion from the created. Through it the Absolute transform itself into all things and then returns back into the emptiness of its undifferentiated nature. In Kaishmiri Shaivism the Absolute is seen as pure consciousness (=being). The Absolute is an eternal all-pervasive principle, the highest reality, the nature of all entities eternally and blissfully at rest within its own nature. The Absolute is the nature of the Self (and

thus of us all). The Absolute is divine; it is Shiva, the Lord of the Universe. It is full of (e) conscious activity through (e&eb)which it generates the universe, and reabsorbs it into (e&eb) itself at the end of each cycle of creation. Thus we speak of monism, as everything resides within (eb) this one absolute consciousness. It sustains (eb) all things, it embraces (e) all things, and it pervades (e&eb) all things. All things are appearances within (eb) the absolute consciousness, but nevertheless real (in contrast to (e) Advaita Vedanta where (e) appearances are seen as unreal or illusionary). All things appear external (out there, outside ourselves), but they do not have (e) a being on their own. They do not exist as (=) separate entities on their own. Everything is contained within (eb) consciousness. What we see as objects are (=) manifestations of consciousness. The event which constitutes (eb) the universe is always internal events happening within (e&eb) consciousness because (e) their essential nature is consciousness itself. If a physical object were totally material, and independent or external to consciousness, it could never be experienced. The universe and (e&eb) consciousness are two aspects of a whole. The universe is (=) an attribute of consciousness which bears consciousness as (=) its substance. Doctrine of Vibration: S.G.Dyczkowski what you "see" is (=) not always what you "get." Many people mistakenly take their own visions literally without (e) "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in (eb) the physical sense but in a psychological sense. The archetypes of (e) the collective unconscious are arrayed behind (e) the scenes of current worldwide conditions, of (e) crisis and confusion. They mirror (e&eb) our own states back at us, whether we perceive them as (=) such or interpret them plausibly or not (Note the essay has been written in model able form and models are given in various papers.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4\,a_9 \square - b_9 \square 0$$

And

$$b_9^{2} \square 4\,a_9 c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9\, in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\, X_9$ ,

$Let\ Y_9 = y_9\, X_9$And try the substitution $A_9 = X_9/(2 + A_9\, y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + n_9 y_9)$ we have

$$W_9^{2} y_9 \left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^{2}\left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9 X_9 \left[\frac{X_9 (1 + n_9 y_9)}{M_9^{2}(2 + n_9 y_9)^{2}} + 2\, y_9 \left(n_9 - 1\right) + \frac{X_9^{2}(9 + y_9)}{M_9^{2}(2 + n_9 y_9)} = 0.\ (A2.12)\right.$$

Since $n_9$ and $A_9$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9 = (2 + n_9 y_9)$, *the condition*(A2.99?)*can be written as*

$[\frac{X_9(E_9-9)}{E_9^2}]^2[\frac{X_9^2}{M_9^2} - 4\, y_9\ E_9^2 \frac{X_9}{M_9^2} + 4y_9 (n_9 - 2)E_9^3]$ .

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9 / W_9^2)$ can be expressed as

$2\, E_9 [y_9 E_9 + (y_9\ E_9 (9 + y_9)(E_9 - n_9))^{9/2}] > \frac{X_9}{M_9^2} > 2E_9 [y_9 E_9 - y_9 E_9 (9 + y_9)(E_9 - n_9))^{9/2}]$

Concept Of Being From Shaivite Point Of View: In His Famous Work Entitled On The Conservation Of Force (1847), Hermann Von Helmholtz Had Divided The Totality Of The Energy In The Universe Between Potential Energy And Kinetic Energy And Affirmed The Reciprocal Convertibility Of The Two. In 1852, William Thomson Pointed Out That There Exists A Sub-Ensemble Within Kinetic Energy, Heat, Which, Once It Has Been Generated, Is No Longer Entirely Convertible Into Potential Energy-Or Into Any Other Form Of Kinetic Energy. Considering That The (Partial) Reconversion Of Heat Into Labor Is Possible Only In Situations That Present A Disparity In Temperature, And That Heat Tends To Pass From Warmer To Cooler Bodies By Spreading On An Even Temperature Level Through Space, Thomson Concluded That The Universe Tends Towards A Final State Where Any Energetic Transformations, Every Movement And Every Form Of Life Will Cease: The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO: The Ultimate Nature Of Anything Cannot Change. Change Of Any Sort Is Merely Apparent (Vivarta). The World Of Change And Becoming Is A False Super- Imposition (Adhyaropa, Adhyasa) On The Absolute. It Appears So On The Screen Of Individual, Collective Consciousness, Albeit One Cannot Be So Sure Of Cosmic, For There Are Many Reasons That Cosmic Transactions Also Had To Be Subjected To Change

By ParaBrahaman. In Cosmic Terms, The Mistake (Bhranti) Consists Of The Supposition That The Real Brahman Is The Unreal Universe And The Unreal Universe Is The Real Brahman. In Microcosmic Terms, It Is The Mistake Of Falsely Conceiving The Body, Mind Or Even One's Personality To Be The Self. Note We Have Taken Self As Witness Consciousness. Notwithstanding There Is Nothing Wrong In The Conception Of Self As A Dynamical Presupposition, In That Individual Consciousness And Collective Consciousness And Even Cosmic Consciousness Evolutes In To ParaBrahaman. In The Same Way As The Image Of A Snake Is Falsely Superimposed On A Rope, Similarly The Universe Is Falsely Projected Onto The Real Substratum, The Brahman. Ignorance Is Not Merely A Personal Lack Of Knowledge, But A Cosmic Principle. As Such It Is Called "Maya," The Indefinable Factor (Anirvacaniya) That Brings This Mistake In Identity About. The Reality Status Of This Cosmic Illusion Is Also Indefinable: On The One Hand It Is Not Brahman, The Sole Reality; On The Other Hand It Is Not Absolutely Non-Existent Like A Hare's Horn Or The Son Of A Barren Woman. Note Here The Vedantin Avoids Talking About Zero Which We Have Avidly And Assiduously Stressed About. Nothing Is The Factor That Exists. Also See the Remarks about Satre's Being And Nothingness. Brahman Is The Source Of World Appearances Only In The Sense Of Being Their Unconditioned Ground Or Essential Nature. The Universe Is False Not Because It Has No Nature Of Its Own But Because It Does Have One. The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski. Concept Of Being From Shaivite Point Of View: It Is The 'I Am The Body/Mind' Belief That Gives Rise To The 'I Am Not The World' Belief. These Two Beliefs Are (=) Co-Created."Consciousness Projects The Appearance Of The Mind, Body And World By Taking (Eb) The Shape Of Thinking, Sensing And Perceiving."“Attention Is Consciousness With An Object. When The Object Vanishes, Attention Simply Remains What It Always Is Consciousness."“There Is No Purpose To Meditation. The Purpose Is Already Accomplished."“Everything

That Is Experienced Is Experienced By, Through, In And As Consciousness.”“The Seen Cannot Be Separated From Seeing And Seeing Cannot Be Separated From Consciousness.” “**The Reality Of Any Experience Is Not Hidden In The Appearance; It Is Expressed By The Appearance The Transparency Of Things**' By Rupert Spira Book Review By Dennis Waite The Following Is A Review Of Rupert's Book: 'The Transparency Of Things: Contemplating The Nature Of Experience’. There Is An Essay From The Book Here And The Essay On Cézanne May Be Read At The Awakened Eye Website Nature Rejects The Naiveté That Seeks Absolute Truth. We Are Beginning To Realize, Individually And Culturally, That "Realities" Are All Human Constructions. The Task Becomes One Of "Catching Ourselves In The Act" Of Creating Our Own "Reality" From The Flow Of Events. Human Truth Is Always An Engagement Of Mind With Experience. The Challenge Of The Therapist In These Times Of Chaotic Change Is To Validate The Concept That We Don't Have To Fear The Collapse Of What We Think We Are. Strange Attractors: Transference, Holography, And An **Archetype Burke, J. (2003). Strange Attractors: Transference, Holography, And An Archetype (Doctoral Dissertation, Pacifica Graduate Institute, 2003).** When we get to the topic and understanding of the Transcendental Logic, the brilliant and baffling Transcendental Deduction (TD) Recall the two movements just discussed, the one from experience to its conditions and the one from the forms of valid inference to the concepts that we must use in all judging (the Categories). This duality led Kant to his famous question of right (quid juris) (A84=B116): with what right do we apply the Categories, which are not acquired from experience, to the contents of experience? (A85=B117). Kant's problem here is not as arcane as it might seem. It reflects an important question: How is it that the world as we experience it conforms to our logic? In briefest form, Kant thought that the trick to showing how it is possible for the Categories to apply to experience is to show that it is necessary that they apply (A97).TD has two sides, though Kant never treats them separately John Marlado,

accounts, attributes, and ascribes the philosophy of Nishida KitarÅ is lucid, insightful, and deeply informative despite the highly questionable argument that gives the book its structure. In comparison with other, technically more accurate accounts of Nishida's philosophy, it is broader in scope than most but more single-minded in its approach. The analysis is for the most part limited to four books that have been translated into English, out of approximately 16 volumes of philosophical essays in Japanese, but it aims to capture the contours and most important details of Nishida's entire philosophy. To appreciate the achievement of this account, one would need only to read, in English translation or in the Japanese original, any of the volumes with their streams of repetitive, puzzling expressions and circuitous, often dead- end argumentation. Wilkinson's elucidation of details in the passages he samples makes it easier to understand why Nishida is considered Japan's greatest academic philosopher. A critical look at the frame of the book and some of its own claims reveals where, to my mind, it has gone astray. The argument that frames this book is repeated throughout: Nishida's philosophy is an attempt to articulate, in a characteristically Western philosophical manner, the Zen experience he had as a young man. Nishida first made the attempt by using the conceptual frameworks of Western philosophers, and after they failed him he developed his own, largely by recasting Buddhist insights. In the end, a philosophical articulation of the Zen worldview in Western terms or even in Eastern terms recast by conceptual, Western philosophical methods approved impossible. What Nishida attempted is incommensurable with Western philosophy. While the author first offers these statements as hypotheses (p. 2), one premise of his argument that there is such a thing as "the Zen world-view" or "the Zen experience “remains unquestioned; and another that Nishida's pivotal experience was the Zen experienceâis asserted as "beyond question" (p. 151; see also pp. 28, 48, 10, 114).Wilkinson's presentation of Zen in the first chapter is drawn largely from D. T. Suzuki, with help from Suzuki's portage Abe Masao; some comments on time, birth, and death in; and a few other sources. The

result is a caricature that imagines Zen as an invariable, unitary experience. In fact, Dagan's teachings on practice as the manifestation of enlightenment do not sit well with D. T. Suzuki's emphasis on a satori experience, and Suzuki does not enlighten us about what Zen in its various practices has historically been. The author would have benefited from acquaintance with other presentations by Zen teachers, such as that of ShunryÅ Suzuki'sZen Mind Beginner's Mind, with historical accounts of Zen traditions, and with contrasts between DÅgen's view of time and Sam Â· sÄraand that of other Buddhist philosophers. A little research in contemporary Zen scholarship would have shown how implausible is the talk of "nirvana or awareness of mu" (p. 24), "the basic Zen position [that] never varies" (p. 3), "the Zen vision" (pp. 5, 7), "Zen epistemology" , "Zen experience" (p. 59), "the conception of time involved in Zen" (p. 15), and "the fundamental Zen assertion that there is no absolute time" (p. 16). Even if DÅgen's notion that time and being are nondual (p. 16) epitomized most of Zen, Nishida's early notion of "a transcendent, unchanging reality apart from time" (p. 51) would contrast sharply with it. Equally questionable is Wilkinson's contention that Nishida's "pure experience" is equivalent to Zen satori. The notion of pure experience Nishida developed in his first major work may be pivotal for his entire philosophy, but he never claims or implies that any part of his philosophy presupposes an exceptional experience like satori that grounds his convictions and renders them unverifiable to those who have not had the experience, as Wilkinson intimates (pp. 55, 159, and 161). Nishida presents his notions of pure experience and intellectual intuition, the grasp of its unity, and later... Perennialism and Constructivism is studied by Randolph T Dible : experiencer or the observer and the transcendental subjectivity itself Pure objectivity is essentially a deconstruction, apophatic and negativa The encountering of the experiencer or observer—transcendental subjectivity itself—at the foundation of the world leads inevitably to the recognition of pure objectivity as ultimate reality (which can be taken as its ultimate deconstruction, analogous to the apophatic or via

negativa), from which objects derive their value, weight, significance, meaning or objectivity. In this way, pure objectivity can be seen as the supra-self-evident Axiological Axiom, so to speak, even Unconditional Love, in romantic terms. This axiology (value theory) has a structure inverse to the relationship between transcendental subjectivity as the radical unity of pure self-reference and on the other hand, the world of forms, as mere traces (representations, indications) of the unique, original "first distinction" Spencer-Brown speaks of at the foundation of his calculus. That is, all forms (i.e., distinctions, differences) would reduce to being the first distinction, also known as the marked state, which can be called penultimate reality (pure self-reference or transcendental subjectivity: the Spirit which animates us), except that forms are complimentary to their content, which is their objectivity or value, which would reduce to the unmarked state or ultimate reality. It is the incongruity of form (thoughts; Whitehead's "negative prehensions") and value (feelings; Whitehead's "positive prehensions," or my notion of objectivity, meaning and qualia; in short, the non-formal aspects of experience) that holds forms open and keeps them from absolute reduction. This accounts for the brute, concrete persistence of the "functional illusion"-- to use a term from Dzogchen Buddhism-- of the world. Thus this system has an axiology of metaphysical objectivity grounded on the ideal of pure objectivity as the source of all value, meaning and significance, itself the very fecundity of profundity, which is the motive of drawing the distinction in the first place **(Wikipedia, Kant's writings and Logic Of Sense)** Nano rings were created by accident while intending to make quantum dots. They have interesting optical properties associated with excitons and the Aharonov–Bohm effect. Application of these rings used as light capacitors or a buffer includes photonic computing and communications technology. Analysis and measurement of geometric phases in mesoscopic rings is ongoing. Several experiments, including some reported in 2012, show Aharonov-Bohm oscillations in charge density wave (CDW) current versus magnetic flux, of dominant period h/2e through CDW rings up to 85 **μm** in

circumference above 77 K. This behavior is similar to that of the superconducting quantum interference devices (seeSQUID). **Aharonov–Bohm nano rings (Wikipedia) Stability analysis of nano rings would help progressively accentuate the research about the materials division that uses nano rings.** From the profound contemplation and wisdom of the Buddha and Mahabodhisattvas, we know there is no such reality. Instead, egolessness (non-self) is the only path to understand the reality of the deluded life. All existences are subject to the law of causes and conditions. These include the smallest particles, the relationship between the particles, the planets, and the relationship between them, up to and including the whole universe! From the smallest particles to the biggest matter, there exists no absolute independent identity. Egolessness (non-self) implies the void characteristics of all existence. Egolessness (non-self) signifies (eb) the non-existence of permanent identity for self and existence (dharma). Sunyata stresses (eb) the voidness characteristic of (e) self and existence (dharma). Sunyata and Egolessness possess (e) similar attributes. As we have discussed before, we can observe the profound significance of Sunyata from (e) the perspective of inter-dependent relationships. Considering dharma-nature and (e&eb) the condition of nirvana, all existences are immaterial and of (e) a void-nature. Then we see each existence as (=) independent of each other. But then we cannot find any material that does exist independent of (e) everything else. So Egolessness also implies (eb) void-nature! From the observation of all existences, we can infer (eb) the theory of nirvana and the complete cessation of (e) all phenomena. From the viewpoint of phenomena, all existences are so different from (e) each other, that they may contradict (e&eb) each other. They are (=) so chaotic. In reality, their existence is (=) illusionary and arises from (e) conditional causation. They seem to exist on one hand, and yet do not (e) exist on the other. They seem to be united, but yet they are so different to (e) one another. They seem to exist and yet they do cease! Ultimately everything will return to (e) harmony and complete calmness. This is the nature of all existence. It is the

final resting place for all. If we can understand this reality and remove our illusions, we can find this state of harmony and complete calmness. **Teachings in Chinese Buddhism (6) sunyata (emptiness) in the Mahayana context (Wikipedia) Many models have been given concatenating Buddhist doctrines and Vaccum energy. These models essentially throw light on the essence of self(witness consciousness) and the cosmic general ledger (cosmic consciousness) while the imbalance in the gratification deprivation complex always remain unbalanced when it comes to individual general ledger , advocating the evolutionary theory of individual consciousness to cosmic consciousness. It is to be noted that it is not the question of knowing all the transactions but the wisdom of what lies behind the debit- Credit, Debit-Debit, or Credit-Credit principle that is most important for assimilation.** Look at this simple thesis to prove the point stated in the foregoing. Thesis: 'The World Has A Beginning In Time, And In Space It Is Also Enclosed In Boundaries.' Proof: 'For If One Assumes That The World Has No Beginning In Time, Then Up To Every Given Moment In Time An Eternity Is Elapsed, And Hence An Infinite Series Of States Of Things In The World, Each Following Another, Has Passed Away. But Now The Infinity Of A Series Consists Precisely In The Fact That It Can Never Be Completed Through A Successive Synthesis. Therefore An Infinitely Elapsed World-Series Is Impossible, So A Beginning Of The World Is A Necessary Condition Of Its Existence, Which Was The First Point To Be Proved (Kant, Critique Of Pure Reason, Tr. By R. Guyer And A. W. Wood (Cambridge: Cambridge University Press, 1998): B 454, P. 470). As Quoted In The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO"The Cyclical Hypothesis, So Heavily Criticized By Nietzsche (VP II 325 And 334), Arises In This Way." In Fact, Nietzsche Was Not Criticizing The Cyclical Hypothesis But Only The Particular Form Of That Hypothesis Presented In Vogt's Work. All Of Nietzsche's Texts Without Exception Speak Of The Eternal Return As The Repetitio Of The Same Events Within A Cycle Which Repeats Itself

Eternally. If Deleuze's Interpretation Holds That The Eternal Return Is Not A Circle, Then What Is It? A Wheel Moving Centrifugally, Operating A "Creative Selection," "Nietzsche's Secret Is That The Eternal Return Is Selective" Says Deleuze: The Eternal Return Produces Becoming-Active. It Is Sufficient To Relate The Will To Nothingness To The Eternal Return In Order To Realize That Reactive Forces Do Not Return. However Far They Go, However Deep The Becoming-Reactive Of Forces, Reactive Forces Will Not Return. The Small, Petty, Reactive Man Will Not Return. Affirmation Alone Returns, This That Can Be Affirmed Alone Returns, Joy Alone Returns. Everything That Can Be Denied, Everything That Is Negation, Is Expelled Due To The Very Movement Of The Eternal Return. We Were Entitled To Dread That The Combinations Of Nihilism And Reactivity Would Eternally Return Too. The Eternal Return Must Be Compared To A Wheel; Yet, The Movement Of The Wheel Is Endowed With Centrifugal Powers That Drive Away The Entire Negative. Because Being Imposes Itself On Becoming, It Expels From Itself Everything That Contradicts Affirmation, All Forms Of Nihilism And Reactivity: Bad Conscience, Ressentiment..., We Shall Witness Them Only Once. [...] The Eternal Return Is The Repetition, But The Repetition That Selects, The Repetition That Saves. Here Is The Marvelous Secret Of A Selective And Liberating Repetition. There Is No Need To Remind The Reader That Neither The Image Of A Centrifugal Movement Nor The Concept Of A Negativity-Rejecting Repetition Appears Anywhere In Nietzsche's Writings, And Indeed Deleuze Does Not Refer To Any Text In Support Of This Interpretation. Further, One Could Highlight That Nietzsche Never Formulates The Opposition Between Active And Reactive Forces, Which Constitutes The Broader Framework Of Deleuze's Interpretation. For Some Years, Marco Brusotti Has Called Attention To The Fact That Deleuze Introduced A Dualism That Does Not Exist In Nietzsche's Writings. To Be Sure, The German Philosopher Describes A Certain Number Of "Reactive" Phenomena (For Example, In The Second Essay Of The Genealogy Of Morality, § 11, He Talks About

"Reactive Affects" [Reaktive Affekte], "Reactive Feelings" [Reaktive Gefühlen], Reactive Men [Reaktive Menschen]); But These Are Nonetheless The Result Of Complex Ensembles Of Configurations Of Centers Of Forces That Remain In Themselves Active. Neither The Word Nor The Concept Of "Reactive Forces" Ever Appears In Nietzsche's Philosophy. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO Some Parts Are Deleted Due To Spatial Constraints. In the midst of thoughts about the eternal return we find at least two other thematic axes On the one hand, the view of the **world as a constant flux of forces without any goal, law, or rules of becoming. A chaos sive natura de-divinized and de-anthropomorphized which constitutes the "ontological substratum**" of the whole of Nietzsche's reflexions On the other hand, an ensemble of fragments of an anthropologico-sociological character, designing a path of liberation leading to the creation of superior individuals by way of a profound transformation of their instinctual structure. This transformation must be achieved by a practice of solitude and internal struggle towards the liberation from the ancient representations of the world and from the incorporated herd values. For an analysis of these thematic perspectives, see Paolo D'Iorio, La linea e il circolo. Cosmologia e filosofia dell'eterno ritorno in Nietzsche (Genova: Pantograf, 1995): 233-322. With most of the parameters quantified, with assumptions there must be picturesque view of the self and non self, space and time, science and philosophy with the concatenational consummations. After The Chapter "On Redemption," Where Zarathustra Dares Not Expose His Doctrine, The Eternal Return Begins To Be Enunciated In Part Three Of The Work. In The First Place, It Is The Dwarf Who Formulates It In The Chapter "On The Vision And The Riddle." Facing The "Gate Of The Instant" Which Symbolizes The Two Infinities That Stretch Towards The Past And The Future, The Dwarf Whispers: "All Truth Is Crooked, Time Itself Is A Circle" The Dwarf Represents The Spirit Of Gravity, And He Embodies The Herd Morality, "The Belittling Virtue" Which Is The Title Of Another Chapter From Part III. The Dwarf Can

Endure The Eternal Return Without Great Difficulties Because He Has No Aspirations; Unlike Zarathustra He Does Not Wish To Climb The Mountains That Symbolize Elevation And Solitude. In Two Unpublished Notes, From The Summer And The Fall Of 1883, Nietzsche Writes: The Doctrine Is At First Favored By The Rabble, Before It Gets To The Superior Men. The Doctrine Of Recurrence Will First Smile To The Rabble, Which Is Cold And Without Any Strong Internal Need. It Is The Most Ordinary Of Life Instincts, Which Gives Its Agreement First. Hence, The Content Of The Doctrine Is The Same, But Whereas The Dwarf Can Endure It (Because He Interprets It According To The Pessimistic Tradition For Which "Nothing Is New Under The Sun"), Zarathustra, Who Is The "Advocate Of Life" Regards The Eternal Return As The Strongest Objection To Existence, And As The Rest Of The Dream Suggests, He Does Not Yet Succeed In Accepting It After The Vision At The Gate Of The Instant, The Chapter Is Brought To An End By The Enigma Of The Shepherd. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO .How does the memory affect the GD complex or for that matter the gratification deprivation of an individual or a collective identity? Non-Markovian local-in-time master equations give a relatively simple way to describe the dynamics of open quantum systems with memory effects. Despite their simple form, there are still many misunderstandings related to the physical applicability and interpretation of these equations. Here, we clarify these issues both in the cases of quantum and classical master equations. Further introduction of the concept of a classical non-Markov chain signified through negative jump rates in the chain configuration. **E-M Laine et al 2012 J. Phys. B: At. Mol. Opt. Phys. 45 154004 doi:10.1088/0953-4075/45/15/154004 Local-in-time master equations with memory effects: applicability and interpretation.** Supposing That There Were Indeed An "Energy Of Contraction" Constant In All Centers Of Force Of The Universe, It Remains To Be Explained Where Any Difference Would Ever Originate. It Would Be Necessary For The Whole To Dissolve Into An Infinite Number Of

Perfectly Identical Existential Rings And Spheres, And We Would Therefore Behold Innumerable And Perfectly Identical Worlds COEXISTING [Nietzsche Underlines This Word Twice] Alongside Each Other. Is It Necessary For Me To Admit This? Is It Necessary To Posit An Eternal Coexistence On Top Of The Eternal Succession Of Identical Worlds? The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. Science is usually portrayed as dehumanizing. Brave New World epitomizes this fear. "The more we understand the world, the more it seems completely pointless" (Steven Weinberg). Certainly science can seem chilling when conceived in the abstract as a metaphysical world-picture. We may seem to find ourselves living in a universe with all the human meaning stripped out: Participants in a soulless dance of molecules, or harmonics of pointlessly waggling superstrings and their braneworld cousins. Nature seems loveless and indifferent to our lives. What rights have we to be happy? **(Brave New World: Aldous Huxley).** Religion can afford to be dogmatic but science cannot. Look at this passage which counters both scientists and philosophers alike: However, Hartmann Objects That The Regressive Movement Postulated By Schopenhauer Is Possible Only In Thought: It Remains Nothing More Than An "Ideal Postulate" With No Real Object And Which "Does Not Teach Us Anything About The Real Process Of The World That Unfolds In A Movement Contrary To This Backwards Movement Of Thought" (Hartmann, Philosophie Des Unbewussten, Third Edition (1871): 772). Hartmann Affirms That If Unlike Schopenhauer One Admits The Reality Of Time And Of The World Process, One Must Also Admit That The Process Must Be Limited In The Past And Therefore That There Must Be An Absolute Beginning. In Hartmann's Mind, Failure To Do So Would Result In Positing The Contradictory Concept Of An Accomplished Infinity: "The Infinity That From The Point Of View Of Regressive Thinking, Remains An Ideal Postulate, Which No Reality May Correspond To, Must, For The World, Whose Process Is, On The Contrary, A Progressive Movement, Open Up To A Determinate Result; And Here The Contradiction Comes To Light" (Hartmann

(1871): 772). What Really "Comes To Light" In This Passage Is The Fact That Hartmann Does Not Provide A Demonstration But A Petitio Principii. Indeed, The Concept Of The World Process Analytically Contains The Concept Of A Beginning Of The World. In All Rigors, It Is Therefore Impossible To Demonstrate These Concepts With Reference To Each Other. Secondly, Hartmann's View That One Is Bound To Accept The Reality Of The World Process Even If One Rejects The Ideality Of Schopenhauer's Time Is Mistaken. Hartmann Believes That If Time Is Real There Must Be A World Process With Both An Absolute Beginning And An Absolute End. Without Any Justification, Hartmann Jumps From Schopenhauer's Negated Time To Oriented Time. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. In the last years several theoretical papers discussed if time can be an emergent property deriving from ) quantum correlations. Here, to provide an insight into how this phenomenon can occur, authors present an experiment that illustrates Page and Wootters' mechanism of "static" time, and Gambini et al. subsequent refinements. A static, entangled state between a clock system and the rest of the universe is perceived as evolving by internal observers that test the correlations between the two subsystems. They implement this mechanism using an entangled state of the polarization of two photons, one of which is used as a clock to gauge the evolution of the second: an "internal" observer that becomes correlated with the clock photon sees the other system evolve, while an "external" observer that only observes global properties of the two photons can prove it is static.Phys. Rev. A 89, 052122 (2014)DOI:10.1103/PhysRevA.89.052122 arXiv: 1310.4691 [quant-ph] **Time from quantum entanglement: an experimental illustration Ekaterina Moreva.** Time is an entanglement .When the new ideas of quantum mechanics spread through science like wildfire in the first half of the 20th century, one of the first things physicists did was to apply them to gravity and general relativity. The results were not pretty. It immediately became clear that these two foundations of modern physics were entirely

incompatible. When physicists attempted to meld the approaches, the resulting equations were bedeviled with infinities making it impossible to make sense of the results. Then in the mid-1960s, there was a breakthrough. The physicists John Wheeler and Bryce DeWitt successfully combined the previously incompatible ideas in a key result that has since become known as the Wheeler-DeWitt equation. This is important because it avoids the troublesome infinites—a huge advance. But it didn't take physicists long to realise that while the Wheeler-DeWitt equation solved one significant problem, it introduced another. The new problem was that time played no role in this equation. In effect, it says that nothing ever happens in the universe, a prediction that is clearly at odds with the observational evidence. This conundrum, which physicists call 'the problem of time', has proved to be thorn in flesh of modern physicists, who have tried to ignore it but with little success. Then in 1983, the theorists Don Page and William Wootters came up with a novel solution based on the quantum phenomenon of entanglement. This is the exotic property in which two quantum particles share the same existence, even though they are physically separated. Entanglement is a deep and powerful link and Page and Wootters showed how it can be used (e) to measure time. Their idea was that the way pair of entangled particles evolves is a kind of clock that can be used to measure change. But the results depend on how the observation is made. One way to do this is to compare the change in the entangled particles with an external clock that is entirely independent of the universe. This is equivalent to god-like observer outside the universe measuring the evolution of the particles using an external clock. In this case, Page and Wootters showed that the particles would appear entirely unchanging—that time would not exist in this scenario. But there is another way to do it that gives a different result. This is for an observer inside the universe to compare the evolution of the particles with the rest of the universe. In this case, the internal observer would see a change and this difference in the evolution of entangled particles compared with everything else is an important a measure of time. This is

an elegant and powerful idea. It suggests that time is an emergent phenomenon that comes about because of the nature of entanglement. And it exists only for observers inside the universe. **Time is an entanglement** (Wikipedia). Danish religious philosopher Soren Kierkegaard carried out a systematic critique of the pretensions of reason and an abstract rationalism which he believed that the modern age was nurturing. Condemning reflection as a "danger" that ensnares people in logical delays and machinations, Kierkegaard compared it to a prison. Reflection is for him a form of captivity, a bondage which "can only be broken by [passionate] religious inwardness" (1978: 81). **Reflection seduces individuals into thinking its possibilities are "much more magnificent than a paltry decision"** (1978: 82). It leads them to act **"on principle,"** to dwell on the deliberation of the context of their actions and the calculation of their worth or outcome. Kierkegaard argues that this drives away feeling, inspiration, and spontaneity, all of which are crucial for true inner being and a vital relation to God. For Kierkegaard, as Nietzsche would later agree, genuine inner being (and culture) is characterized by the tautness and tension of the soul which characterizes passionate existence. But the "coiled springs of life relationships ... lose their resilience" in reflection (1978: 78) and "everything becomes meaningless externality, devoid of [internal] character" (1978: 62). Kierkegaard thus contributes to the development of an irrationalist tradition that has echoes in some later postmodern thought. Kierkegaard might well have agreed with his contemporary Fyodor Dostoyevsky, who wrote: "An intelligent [reflective] man cannot seriously become anything ... excessive consciousness is a disease" (1974: 3, 5). In an age overtaken by rules and regulations, genuine action -- which Kierkegaard assumes to be subjective and spontaneous -- is frustrated at every turn. Complaining that we are too "sober and serious" (1978: 71) even at banquets, Kierkegaard bemoans the fact that even suicides are premeditated (1978: 68)! "That a person stands or falls on his actions is becoming obsolete; instead, everybody sits around and does a brilliant job of bungling through with the aid of some reflection

and also by declaring that they all know very well what has to be done" (1978: 73). Thus, it is passion, not reflection, that guarantees "a decent modesty between man and man [and] prevents crude aggressiveness" (1978: 62). "Take away the passion and the propriety also disappears" (1978: 64). The ambiguity in the word "passion" may cause some confusion here. To say that the age and its individuals are "passionless" is not to say there are no emotions whatsoever, but rather that there is no true spiritual inwardness and depth, no intensively motivated action and commitment. It suggests that passion exists only in a simulated, pseudo-form, "the rebirth of passion" through "talkativeness" (1978: 64). "Chattering" for Kierkegaard gets in the way of "essential speaking" and merely "reflects" inconsequential events (1978: 89-99). Hence, in "the present age," emotions -- which in fact are all too pronounced -- have been transformed into negative forces. Anticipating Nietzsche's genealogy of the "slave revolt" in morality, Kierkegaard claims that the "enthusiasm" of the prior age of Revolution, a "positively unifying principle," has become a vicious "envy," a "negatively unifying principle" (1978: 81), a leveling force in its own right insofar as those lacking in talent and resources want to tear down those who have them. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS By Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html.** In November 2011, the United States government argued before the US Supreme Court that it wants to continue utilizing GPS tracking of individuals without first seeking a warrant. In response, Justice Stephen Breyer questioned what this means for a democratic society by referencing Nineteen Eighty-Four. Justice Breyer asked, "If you win this case, then there is nothing to prevent the police or the government from monitoring 24 hours a day the public movement of every citizen of the United States. So if you win, you suddenly produce what sounds like 1984... "[65]. What is to be said in unmistakable terms is that when it comes to national security and innocent citizens there is question of compromise even it be an infringement and encroachment of so called

privacy which has been made a word to subterfuge the activities that are far said to be fair. Quantum theory has provoked intense discussions about its interpretation since its pioneer days. One of the few scientists who have been continuously engaged in this development from both physical and philosophical perspectives is Carl Friedrich von Weizsäcker. The questions he posed were and are inspiring for many, including the authors of this contribution. Weizsäcker developed Bohr's view of quantum theory as a theory of knowledge

## ABSOLUTE STABILITY

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_9$ ) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9 - 9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2 - 2y_9\frac{(E_9 + n_9^2 y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right) - 4y_9(E_9 - n_9)$$

And has the solution

$$\frac{X_9}{M_9^2} > \frac{E_9}{(E_9 - 9)}[y_9(E_9 + n_9^2 y_9) + [y_9^2 (E_9 + n_9^2 y_9)^2 + 4y_9(E_9 - 9)(E_9 - n_9)]^{9/2}]$$

Given $E_9 > n_9$this lower limit to $X_9 / M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9 / M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9 < n_9$ $(ie., 2A_9 M_9^2 < X_9 - Y_9$)then (A2.9)and (A2.8)show that $Q_9$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2} > 2E_9[y_9 E_9 - (y_9 E_9 (9 + y_9)(E_9 - n_9)]^{9/2}]$$

And

$$E_9 > n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2/2M_9 (X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$and $X_9$and $Y_9$by $W_9^2$ the dependence upon *M* disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9$ *and* separate (A3.1)into real and imaginary parts

$$Z_9^2 + (M_9 + V_9) Z_9 \sin(Z_9 T_9) - (U_9 M_9 + Y_9) + (M_9(U_9 + V_9) - X_9)$$

$$\mathrm{X}\cos(Z_9 T_9) - 2M_9 V_9 \cos^2(Z_9 T_9) + M_9 V_9 = 0$$

$$(M_9 + U_9) Z_9 - (M_9 + V_9) Z_9 \cos(Z_9 T_9) + (M_9(U_9 + V_9) - X)$$

$$\mathrm{X}\sin((Z_9 T_9) - M_9 V_9 \sin(Z_9 T_9) = 0$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 = 1 - \cos Z_9 T_9$and $A_9 = U_9 - V_9$we obtain

$$4 M_9^2 V_9^2 W_9^4 + 2 M_9 V_9 W_9^3[(M_9 - V_9)^2 + A_9(M_9 + V_9) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2[M_9 - V_9)^2(2 A_9 M_9 + Y_9 - X_9) + (2A_9 M_9 - X_9)^2 - A_9(M_9 - V_9)(2A_9 M_9 - X_9)$$

$$+ M_9 V_9 (A_9^2 + 4(X_9 + Y_9))]$$

$$+ W_9[A_9^2 + 2X_9)(2 A_9 M_9 - X_9) + 2A_9 Y_9(M_9 + V_9) - M_9 A_9^3] + A_9^2(X_9 M_9 - Y_9) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9(2, U_9 V_9) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting cos $(Z_9 T_9) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MDDULE NUMBERED SEVEN**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability

criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

## Morphological And Epistemological And Syllogistic Models For Which Nunny's Model Could Be Applied And Has Been Applied:

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production**

**Charge Asymmetry, $\Delta Y \equiv Yt - Yt^{-}$, Reflects The Asymmetry In $Tt^{-}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles

ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic

quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation

between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to

nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of

Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first

impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of

another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease. Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation

of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces

always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the

thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of

thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity

and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI

Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs” Outcomes—

And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities

Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \quad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

## DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t-T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\big(R(t)\big)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and haimovici model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

**TIME DELAY**

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Limitations of Predator prey analysis in immense. When we realize that world itself is Debit-Credit world and sometime credit-Credit world (See an earlier paper entitled the same nomeneclature.

Opposite limit of absolute stability. The result is a specific The Doctrine of Vibration awareness of an object of the form: 'here this is such.' While individual appearances do not lose their separate identity {svarupabheda) when they rest on a common basis, even so the particular object which appears according to its own characteristics (svalak$ana) is an individual reality in its own right. It is a different kind of appearance characterised by its association with the appearance of the specific location and time in which it is made manifest. The form of our experience is thus 'I now see this here'. But when we perceive each particular constituent appearance separately, each assumes a separate fixed function. Abhinava cites the following colourful example to illustrate how the various combinations of appearances account for the variety of experience: Thus even though the appearance of the beloved may manifest externally, it is as if far away in the absence of another appearance, namely, that of 'embracing'. So when the [appearing of the beloved] is associated with another appearance [namely that of 'far away'] the power (arthakriya) it formerly had of giving pleasure appears as its contrary. The form our experience assumes depends, not only on the nature of the object perceived, but also on personal factors entirely peculiar to ourselves. This theory explains this in two ways. In one sense, the object remains the same, but one or other of its constituent appearances

comes to the fore according to the inclinations of the perceiver. From another point of view, we can say that the perceived object is different for each perceiver according to the difference in the prominent appearance manifest to him. Abhinava, citing as an example a golden jar, illustrates how the same object appears differently to different perceivers according to the use they wish to make of it and to their state of mind: When a person who is depressed and feels that there is nothing [of value for him in the world] sees the jar, he merely perceives the appearance 'exists' [in the form of the awareness that] 'it is\ He is not conscious of any other [of its constituent appearances] at all. An individual who desires to fetch water [perceives] the appearance 'jar'. The man who simply wants something that can be taken somewhere and then brought back [perceives] the appearance 'thing'. The man who desires money [perceives] the appearance 'gold'. The man who desires a pleasing object [perceives] the appearance 'brightness' while he who wants something solid sees the appearance 'hardness'. These 'atomic events' or appearances emerge from the pure subject's consciousness and combine together to form a total event at each moment. Integral Monism of Kashmiri Saivism Daily life (vyavahara) goes on by virtue of this ever renewed flux of appearances. They are connected together and work towards a single unified experience because they appear within the field of consciousness of the universal subject. The aggregate of appearances arises in the [supreme] subject as do [sprouts in] a rice field. Even though each sprout germinates from its own seed, they are perceived as a collective whole. Appearances rest in this way within the universal subject. 'External- ity' is itself another appearance; it arises from a distinction between appearances and the individual subject. So, although all manifestation always occurs within the subject, it appears to be external due to the power of Maya which separates the individual subject from his object. This split must occur for daily life to be possible. Only externally manifest appearances can perform their functions; when they are merged within the subject and at one with him, they cannot do so. Daily life proceeds on the basis of the operation and withdrawal of the conditions necessary for fruitful action to be possible. Appearance in this sense represents the actualisation of a potential hidden in consciousness made possible by virtue of its dynamic, Spanda

nature which is both the flow from inner to outer and back as well as the power that impels it. The emergence from, and submergence into, pure consciousness of each individual appearance is a particular pulsation (visesaspanda) of differentiated awareness. Together these individual pulsations constitute the universal pulse (samanyaspandd) of cosmic creation and destruction. Thus, every single thing in this way forms a part of the radiant vibration {sphuratta, sphurana) of the light of absolute consciousness. Light and Awareness: The Two Aspects of Consciousness Absolute consciousness understood as the unchanging ontological ground of all appearing is termed 'Prakasa'. As the creative awareness of its own Being, the absolute is called ' Vimarsa '. Prakasa and Vimarsa — the Divine Light of consciousness and the reflective awareness this Light has of its own nature — together constitute the all-embracing fullness (purnata) of consciousness. The Recognition (pratyabhijfia) school of Kashmiri Saivism develops this concept of the absolute which finds its fullest expression in Utpaladeva's Stanzas on the Recognition of God. Even though neither of these two key terms appear in the Stanzas on Vibration or the Aphorisms of Siva, they recur frequently in their commentaries. Thus, although the original formulation of the Doctrine of Vibration differs from the theology of Recognition in this respect, it was extended in the course of its development to accommodate this concept of the absolute as well. This was possible, and quite justified, insofar as the absolute understood in Pratyabhijfia terms does not, as we shall see, differ essentially from that of the Spanda school. We can, as Kashmiri Saivites themselves have done, explain one in terms of the other. The Doctrine of Vibration Prakasa: The Light of Consciousness Prakasa is the pure luminosity' (bhdna) or 'self-showing' that constitutes the essence and ultimate identity (atman) of phenomena. That things appear at all is due to the light of consciousness, and their appearing (avabhasana) is itself this Light which bestows on all things their evident, manifest nature. Established in the light of consciousness everything appears there according to its own specific nature (svabhava). Anything that supposedly does not rest in this Light is as unreal as a sky-flower. 3 Thus, according to Rajanaka Rama, unlike the light of the sun, or any other light, this Light not only makes all things apparent, it is also their ultimate source. 4 Full of its

divine vibration the Light makes all things manifest and withdraws them into itself. This supra- temporal activity characterises it most specifically; devoid of it, it would be no better than an inert physical phenomenon. At the same time, this light is the conjunction (slesa) or oneness (aikdtmya) of its countless manifest forms, 6 and the collective whole (Sampina'ana) of all the categories of existence. The universe is nothing but the shining of the Light within itself. It is the radiant vibration (sphuratta) of this Light, the state (avast hand) in which consciousness becomes manifest. Although the Light shines as all things at all times and hence also makes their diversity manifest, 9 penetrating each object individually as well as collectively, it is not totally 'merged' (magna) or identified with the object so as to suffer any division within itself. Our experience of any object is of the form: T see this': it is not itself an object, but the manifest form the object assumes as a luminous principle of experience. The Light is ever revealed and can never be obscured; objectivity can never cast a shadow on the light of consciousness. The Stanzas on Vibration declare: That in which all this creation is established and from whence it arises is nowhere obstructed because it is unconditioned by [its very] nature. This Light is the highest reality (Paramartha). lity with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**DELAY DUE TO CONFEDERATION DEBIT TYPE MATURATION CREDIT TYPE CONSOLIDATION MATURATION**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \qquad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9} + \frac{N_9 dG_9}{dR_9} - \frac{dB_9}{dR_9} > \frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9} + \frac{dM_9}{dR_9}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened**

**eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production Charge Asymmetry, $\Delta Y \equiv Yt - Y\bar{t}$, Reflects The Asymmetry In $T\bar{t}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton

particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results

in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely

dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact**

**most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the**

**thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is

nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the

world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease.

Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition

and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the

second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular

chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon

be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining

Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been

Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing

Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

## ABSOLUTE STABILITY

$$L_9 = \frac{N_9 G_9}{2M_9} \frac{d(9nF_9)}{dR_9} (\geq 0) \quad (9a)$$

$$I_9 = -\frac{N_9 G_9}{2M_9} \frac{d(9nM_9}{dR_9} (\geq 0) \quad (9b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9 > L_9 - I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_9 > L_9 - I_9$, when $P_9^2/4 > I$ (8a)

$P_9 > L_9 - I_9 + \frac{(I_9 - P_9^2}{(L_9 - I)}$ when $P_9^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to

reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

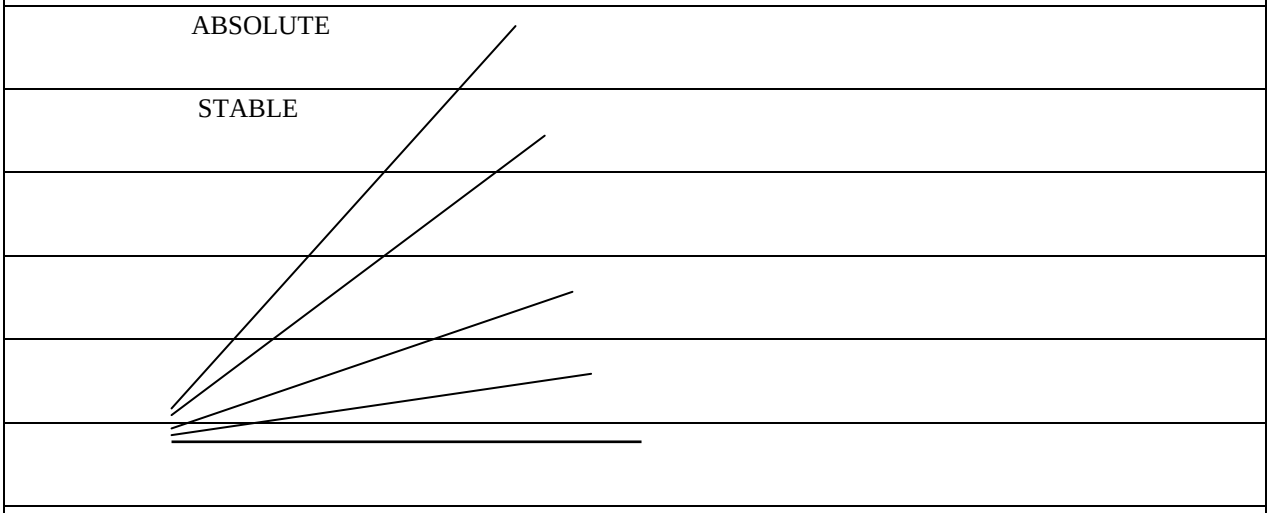



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori

reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

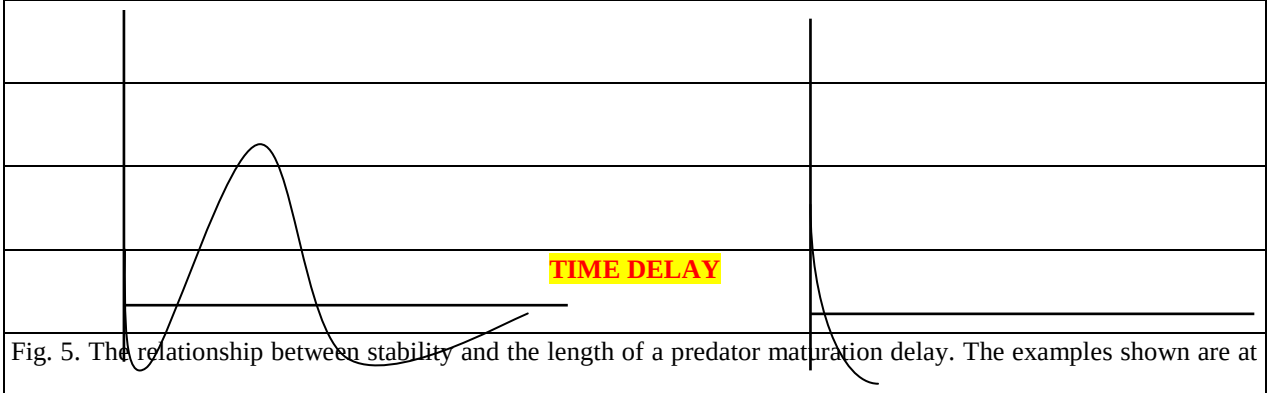



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

**DELAY IN BOTH PREDATOR AND PREY VIS A VIS DEBIT TYPES AND CREDIT TYPES**

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t - T_9)F_9\big(R_9(t - T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t - T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ and the effective prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting

$dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

Systems to which model has been applied also include: **Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI:

http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors.

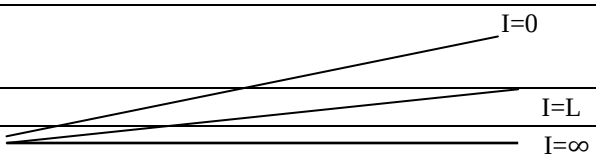
Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, ,**

**Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)**

TOTAL

ABSOLUTELY STABLE

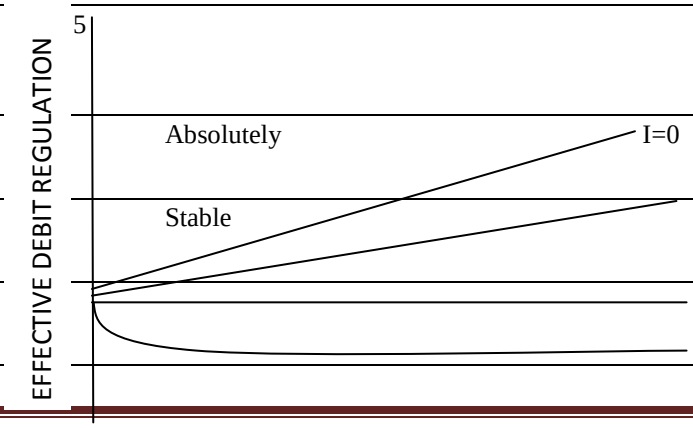



DELAYED DEBIT EFFECT

FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.



-50 10

**DELAYED CREDIT EFFECT (LACK OF BALANCE) CREDIT NON CUMSUMMATION:"**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

The following treatise bears ample testimony and infallible observatory to the fact that how Nunny's model can be applied to various cosmological, general relativity and economic problems.

## ONTOLOGICAL AND EPISTEMOLOGICAL FEATURES OF THE STUDY: OF NUNNY's MODEL

When they are known to be as they are, in the fullest sense of their existence, their phenomenal nature disappears leaving the ground of Being naked and accessible. This approach was validated by a critique of experience. The Vedanta established that space, time and the other primary categories of our daily experience can have no absolute existence. It was therefore necessary to make a distinction between relative truths — that accepted by the precritical common man — and an absolute truth discovered at a higher level of consciousness. The Saiva absolutist 17 rejects any theory that maintains that the universe is less than real. From his point of view a doctrine of two truths, one absolute and the other relative, endangers the very foundation of monism. The Kashmiri Saiva approach is integral: everything is given a place in the economy of the whole. It is equally wrong to say that reality is either one or diverse. Those who do so fail to grasp the true nature of things which is neither as well as both. "We do not" says Abhinavagupta, "base our contention that [reality] is one because of the contradictions inherent in saying that it is dual. It is your approach {paksa) that accepts this [method] While, if [duality and oneness] were in fact [to contradict each other], they would clearly be two [distinct realities]." The Vedantin, who maintains that non-duality is the true nature of the absolute by rejecting duality as only provisionally real, is ultimately landed in a dualism between the real and illusory by the foolishness of his own excessive sophistry (vacafadurvidya). Oneness is better understood as the coextensive unity (ekarasa) of both duality and unity. They are equally expressions of the

absolute. Gopinath Kaviraj says: According to Sarikara, Brahman is truth and Maya is inexplicable (anirvacaniya). Hence the [Advaitin's] endeavour to demonstrate the superiority of Advaita philosophy is turned against his own system. It tarnishes the picture of its philosophical perfection and profundity. He cannot accept Maya to be a reality, therefore his non-dualism is exclusive. The whole system is based on renunciation and elimination and thus is not all-embracing.... By accepting Maya to be Brahman (brahmamayi), eternal (nitya) and real (satyarupa), Brahman and Maya [in the Tantra] become one and coextensive. **The Doctrine of Vibration: An Analysis of the Doctrines and Practices of Kashmir Shaivism Mark S. G. Dyczkowski. Concept of being from Shaivite point of view:** Solutional behaviour and stability analysis of such systems is of paramount importance, notwithstanding the assumptions made which are not farfetched, in that it throws light on the fact whether they could be candidates for yet unsolved problems, providing Autheticative determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication for the wide amplitudinal ramificatory usage of the thesis propounded. The world is now the other person's world, a foreign world that no longer comes from the self, but from the other. The other person is a "threat to the order and arrangement of your whole world...Your world is suddenly haunted by the Other's values, over which you have no control"(Satre: being and Nothingness) Sartre states that many relationships are created by people's attraction not to another person, but rather how that person makes them feel about themselves by how they look at them. This is a state of emotional alienation whereby a person avoids experiencing **their subjectivity by identifying themselves with "the look" of the other**. The consequence is **conflict**. In order to maintain the person's own being, the person must control the other, but must also control the **freedom of the other "as freedom**". These relationships are a profound manifestation of "bad faith" as for-itself is replaced with the other's freedom. The purpose of either participant is not to exist, but to maintain the other participant's looking at them. This system is often mistakenly called "love", but it is, in fact, nothing more than emotional alienation and denial of freedom. In

essence we are always simulating and selling ourselves at the cost of our freedom. Both as a script writer and participant consciousness, human beings experienceâ complete subjective feelings and sometimes absolute subjective feelings, which if remains unchecked might lead to pathological ramifications. (Italics mine) Perspicuous forbearance, Sophisticated seasoning, Partcipational observation, Background actuality, Existential worldliness, Existential strife, Predicational anteriorities, Character consonance, Ontological consonance, Primordial exactitude, Phenomenological correlates, Accolytish representation, Atrophied asseveration. Anamensial alienisms, Anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orinetatilities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state form the bastion, pillar, post, stylobate and sentinel of the classification scheme and doxa. appropriate and diagnostic, differentiating-discriminating, distinctive-distinguishing, emblematic- especial, essentially exclusive, fixed, idiosyncratic-inborn, inbred, indicative and individual, individualistic-inherent, individualizing and ingrained, inherent, innate, local, marked, native, normal, original, particularly- peculiar, personal, private, proper, regular, representative-singular, special, specific, symbolic, symptomatic- unique these form the classification and stratification measure that is suitable notwithstanding the constancy in many a case, as in gravity, which depends on the mass and distance for two given objects**.** Like Kant and Bergson, Deleuze considers traditional notions of space and time as unifying forms imposed by the subject. He therefore concludes that pure difference is non-spatio-temporal; it is an idea, what Deleuze calls "the virtual". (The coinage refers to Proust's definition of what is constant in both the past and the present: "real without being actual, ideal without being abstract.") Proust, Le Temps Retrouvé, ch. III: see the fourth line from the bottom of this page, or, in English translation, the thirteenth paragraph here: "I began to discover the cause by comparing those varying happy impressions which had the common quality of being felt simultaneously at the actual moment and at a distance in time, because of which common quality the noise of the spoon upon the plate, the unevenness of the paving-stones, the taste of the madeleine, imposed the past

upon the present and made me hesitate as to which time I was existing in. Of a truth, the being within me which sensed this impression, sensed what it had in common in former days and now, sensed its extra-temporal character, a being which only appeared when through the medium of the identity of present and past, it found itself in the only setting in which it could exist and enjoy the essence of things, that is, outside Time. [...] **Nothing but a moment of the past Much more perhaps; something which being common to the past and the present, is more essential than both. [...] a marvellous expedient of nature had caused a sensation to flash to me—sound of a spoon and of a hammer, uneven paving-stones**—simultaneously in the past which permitted my imagination to grasp it and in the present in which the shock to my senses caused by the noise had effected a contact between the dreams of the imagination and that of which they are habitually deprived, namely, the idea of existence—and thanks to that stratagem had permitted that being within me to secure, to isolate and to render static for the duration of a lightning flash that which it can never wholly grasp, a fraction of Time in its pure essence. When, with such a shudder of happiness, I heard the sound common, at once, to the spoon touching the plate, to the hammer striking the wheel, to the unevenness of the paving-stones in the courtyard of the Guermantes' mansion and the Baptistry of St. Mark's, it was because that being within me can only be nourished on the essence of things and finds in them alone its subsistence and its delight. It languishes in the observation by the senses of the present sterilised by the intelligence awaiting a future constructed by the will out of fragments of the past and the present from which it removes still more reality, keeping that only which serves the narrow human aim of utilitarian purposes. But let a sound, a scent already heard and breathed in the past be heard and breathed anew, simultaneously in the present and in the past, real without being actual, ideal without being abstract, then instantly the permanent and characteristic essence hidden in things is freed and our true being which has for long seemed dead but was not so in other ways awakes and revives, thanks to this celestial nourishment "While Deleuze's virtual ideas superficially resemble Plato's forms and Kant's ideas of pure reason, they are not originals or models, nor do they transcend possible experience; instead

they are the conditions of actual experience, the internal difference in itself. "The concept they [the conditions] form is identical to its object."[Desert Islands, p. 36.] A Deleuzean idea or concept of difference is therefore not a wraith-like abstraction of an experienced thing, it is a real system of differential relations that creates actual spaces, times, and sensations. [See "The Method of Dramatization" in Desert Islands, and "Actual and Virtual" in Dialogues. In fact it not as if the style and substance of science and philosophers have been entirely different. look at the conformality and consonance in the Gaussian Theory and Deleuze's statements on identity and the loss of it in the following passage…….That is, identity is a continuous process that finds its limit in exhaustion. As Deleuze will write, "…the logical relation of causality is inseparable from a physical process of signalling, without which it would not be translated into action. By 'signal' we mean a system with orders of disparate size, endowed with elements of dissymmetry; by 'sign' we mean what happens within such a system, what flashes across the intervals when a communication takes place between disparates. The sign is indeed an effect, but an effect with two aspects: in one of these it expresses, qua sign, the productive dissymmetry; in the other it tends to cancel it" (dr, 20). An actualized entity just is such a signal-sign system, remaining in perpetual communication with the virtual multiplicity it actualizes, constituting its own elements and identity. Such is Deleuze's account of what Badiou refers to as the operation of the "count-for-one"… an operation that is immanent to being itself and energetic in character. yet how are we to discern the local nature of ontological situations, the manner in which there is no global, overarching global situation in which all local situations are embedded as if parts of a whole? in part the above passage already responds to this question. if series must be brought into relation to communicate in order for causality to occur, and if the resonance of series is always a matter of chance, of a throw of the dice, like Lacan's objet a where no series enjoys primacy of model over copy, then all we have are local situations without a total situation. Unlike Heidegger's analysis of Dasein where Dasein is always being-in-the-world, Deleuze's simulacra or actualizations are always divergent, enjoying only local relations. Heidegger argues that in order to

understand the being of Dasein we must understand the manner in which Dasein is a being-in-the-world. By contrast, Deleuze's concept of multiplicity provides us with the principle of a new structuralism, of a new local ontology, that allows us to understand the immanent organization of a multiplicity without referring it to an embedding global space. as DeLanda so beautifully puts it in relation to the mathematician gauss, "…when gauss began to tap into these differential resources, a curved two-dimension surface was studied using the old Cartesian method: the surface was embedded in three-dimensional space complete with its own fixed set of axes; then, using those axes, coordinates would be assigned to every point of the surface; finally, the geometric links between points determining the form of the surface would be expressed as algebraic relations between the numbers. but gauss realized that the calculus, focusing as it does on infinitesimal points on the surface itself (that is, operating entirely with local information), allowed the study of the surface without any reference to a global embedding space. basically, gauss developed a method to implant the coordinate axes on the surface itself (that is, a method to implant the coordinate axes on the surface itself (that is, a method of 'coordinatizing' the surface) and, once points had been so translated into numbers, to use differential (not algebraic) equations to characterize their relations. as the mathematician and historian Morris Kline observes, by getting rid of the global embedding space and dealing with the surface through its own local properties, 'gauss advanced the totally new concept that a surface is a space in itself'" (intensive science and virtual philosophy, 11-12).according to the old Cartesian method, we can only outline the properties of a space by relativizing it to a global space in terms of which it is then mapped. by contrast, gauss is able to explore a space in terms of its intrinsic metric and organization as a local space, without referencing it to a whole of which it is conceptualized as a part. **The Image of Thought** "Image of thought "permeates both popular and philosophical discourse. According to this image, thinking naturally gravitates towards truth. Thought is divided easily into categories of **truth and error**. The model for thought comes from the educational institution, in which a master sets a problem and the pupil produces a solution which is either true or false. This image of the subject

supposes that there are different faculties, each of which ideally grasps the particular domain of reality to which it is most suited. In philosophy, this conception results in discourses predicated on the argument that "**Everybody knows..."** the truth of some basic idea. Descartes, for example, appeals to the idea that everyone can at least think and therefore exists. Deleuze points out that philosophy of this type attempts to eliminate all objective presuppositions while maintaining subjective ones. Deleuze maintains, with Artaud, that real thinking is one of the most difficult challenges there is. Thinking requires a confrontation with stupidity, the state of being formlessly human without engaging any real problems. One discovers that the real path to truth is through the production of sense: the creation of a texture for thought that relates it to its object. Sense is the membrane that relates thought to its other. Accordingly, learning is not the memorization of facts but the coordination of thought with a reality. "**As a result, 'learning' always takes place in and through the unconscious, thereby establishing the bond of a profound complicity between nature and mind".** Deleuze's alternate image of thought is based on difference, which creates a dynamism that traverses individual faculties and conceptions. This thought is fundamentally energetic and signifying: if it produces propositions, these are wholly secondary to its development. Philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; even tum tan tum;it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the vorticity of being merges with the positive use of the disjunctive synthesis, and

this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoin der for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for every hum of voices or a single silence for all the silences; a single vacuum for all the vacuousness consciousness should not be said without conjuring if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. The attributable to all states of bodies and states of affairs and the expressible of every proposition So univocity of consciousness means the identity of thé noematic attribute and that which is expressed linguistically and sense-fully. Univovity.We mean that it does not allow consciousness to be subsisting in aquasi state and but expresses in all pervading reality. (Reference: LOS: Deleuze) Deleuze sometimes spoke of getting up behind an author and "creating a monster", as if his reading methodology somehow consisted of a distortion of the philosopher's or author's thought. Certainly there is a case to be made for this with regard to the process of deterritorialization, where something is wrested from a territory and deterritorialized upon a new territory, like the animal paw that is deterritorialized from the earth and reterritorialize on the branch. however, it seems to me that the more interesting aspect of Deleuze's approach to other philosophers and art is not so much his "monstrous creations" (i seldom find them particularly monstrous), but rather their Gaussian or Riemannian style, where he explores them in terms of their own internal organization and metric, without reducing them to something alien such as history, society, biology (reductivism), or the signifier. What we find in Deleuze's approach to phenomena is a Gaussian technique. for instance, take Deleuze's books on cinema, his book on Francis bacon, or his study of sacher-masoch. In the first instance, Deleuze's carefully separates cinema from narrative and the signifier, studying it in terms of its specific organization pertaining to the production of images. in the case of bacon, Deleuze doesn't look for an underlying narrative or "meaning", but instead studies the manner in which bacon composes and

organizes his images and lines, both in terms of their production and actuality, so as to liberate a logic of sensation. finally, in approaching sacher-masoch, Deleuze allows Sacher-Masoch's novels to speak for themselves in terms of their desire and relation to pain, stalwartly refusing to reduce masochism to the complement of sadism. in each case, we have a local exploration of a "space" of multiplicities that is extremely precise. Sean Bowden the priority of events: **Deleuze's logic of sense sean Bowden, the priority of events: Deleuze's logic of sense, Edinburgh university press, 2011, 296pp., $40.00 (hbk), isbn 9780748643646 Reviewed by David Scott, Coppin state university and Gaussian spaces and multiplicities a note on Gaussian spaces and multiplicities posted by larval subjects under uncategorized.** What you "see" is not always what you "get." Many people mistakenly take their own visions literally without "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in the physical sense but in a psychological sense. The archetypes of the collective unconscious are arrayed behind the scenes of current worldwide conditions, of crisis and confusion. They mirror our own states back at us, whether we perceive them as such or interpret them plausibly or not. The noise of ordinary consciousness and beliefs drowns out the signal. Unconsciousness is the background of our ordinary awareness. Our organism is very much at the center of such effects. The organismic source is our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is (=) a meme-Scape. Pre-conceived concepts vie with (e&eb) structures, concepts with images. Like scientists who ignore (e) assumed truths, we leapfrog over (e&eb) our beliefs and personality deficits, claiming (eb) idiosyncratic imagination is (=) literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves (e) us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our

experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images. It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is (=) both. We inherit (e&eb) the structures that make our experience (=) what it is. But the structure itself is "empty," and each human culture "fills" it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about (e&eb) our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from

criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns,

nomenclature, and internal coherence of meaning and representation. n a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar some statements in Vedanta, it is not superfluous or redundant. It is evolutionary. Authetication determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication of the evolutionary process towards understanding is ever present in all, its thematic and discursive form. it is the relentless holding on to a theory this is despicable albeit the theory is not. Look at this paragraph. It is the 'I am the body/mind' belief that gives rise to the 'I am not (e) the world' belief. These two beliefs are (=) co-created."Consciousness projects (e&eb) the appearance of the mind, body and world by taking the shape of thinking, sensing and perceiving." "Attention is Consciousness with an object. When the object vanishes, attention simply remains what it always is, Consciousness.""There is (=) no purpose to (e) meditation. The purpose is already accomplished." "Everything that is experienced is experienced by, through, in and as Consciousness." "The seen cannot be separated from seeing and seeing cannot be separated from (e) Consciousness.""The Reality of any experience is not hidden in the appearance; it is expressed by the appearance." "Once we see that everything is Consciousness... Maya still dances, but it is a dance of love not seduction." There are some truly excellent sections

such as a long one taking, as a starting point, the following quotation from the painter Cézanne: "Everything vanishes, falls apart, doesn't it? Nature is always the same but nothing in her that appears to us lasts. Our art must render the thrill of her performance, along with her elements, the appearance of all her changes. It must give us a taste of her Eternity." Rupert begins: "That statement must be one of the clearest and most profound expressions of the nature and purpose of art in our era." And he goes on for some 14 pages to elaborate on this claim, examining the nature of the 'elements' Existence, appearance and Consciousness and their relationship

Some of the following concepts are modeled; some are not. For those which are not modelled kindly see one of the papers in the series. This note has the axiomatic predication to give a holistic view of the study. Here we talk of the characterstics of systems which satisfy the condition of cosmological constant or any other for that matter or variable or axiomatic predication. . There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit,

facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified

from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. Towards the end of circumvention of typing hundreds of superscripts and subscripts which would be a sardine squash and the concomitant operational difficulties, model is presented in piece meal of nine modules. Logarithms are to be taken in respect of those which incorporate more than one variable in bra-ket. Values of log (ab) and log (a+b) are readily available. Anti logarithm shall be taken at the end to predict, project, prognosticate the value of the variable. This is true for tensors, vectors and other variables too. Affirmational assertion, and explanation justification, statement of vindication, annotational commentary, and explicational glossary for each and every system changes and physical interpretation of results is one thing that is to be with earnest endeavour and feverish and febrile expediency. Editions never ever mean the same and identicality of parameters and this has been explicitly and unmistakably stated in the model apriori itself. Akin and analogous, cognate and concurrent, correspondent and congruent, comparable and complementary, synonymous and duple, tantamount and agnate, commensurate and correlative representation is only to highlight the importance and subterfuge the replication of work. It is my fervent solicitation to kindly bear with me for any lapses, notwithstanding the orchestrated efforts for a paper without any minor errors. Postulation predication, conclusive presumption, differential presuppositions, underscored decidedly axiomatic statement of the statement, equation form the bye word or the watch word in the aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement-argument, articulation, assertion, asseveration, choice of words, commentary and communication, declaration and definition, delivery-diction, elucidation-emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation tone and tenor of stratification. Now, how do find the reaction of systems to these singularities. You do the same thing a boss does for you. "Problematize" the events and see how you behave. I will resort to "pressure tactics". "intimidation of deriding report", or "cut in the increment" to make you undergo trials, travails and tribulations. I am happy to see if you improve your work; but may or may not be sad if you succumb to it and hang yourself! We do the same thing with systems. systems show conducive response, felicitous reciprocation or behave erratically with inner roil, eponymous radicality without and with blitzy conviction say like a solipsist nature of bellicose and blustering particles, or for that matter coruscation, trepidiational motion in fluid flows, or seemingly perfidious incendiaries in gormandizing fellow elementary particles, abnormal ebullitions, surcharges calumniations and unwarranted(you think so but the system idoes not!) unrighteous fulminations. Amount of emotion or affect that is appropriate or expected in an given form of interaction. Particular individuals and diffuse obligations (see c and d) are associated with affectivity, whereas contacts with many individuals (universalistic) in a bureaucracy may be devoid of emotion and characterized by affective neutrality. Affective neutrality may refer to self discipline and the deferment of gratification (eg Weber's spirit of capitalism). In contrast, affectivity may be associated with expressing emotions. Adams and Sydie also refer to affective neutrality being associated with ego control (p. 15). (Parsons: Wikipedia). So the point that is made here is "like we problematize the "events" to understand the human behaviour we have to "problematize" the events of systems to understand their behaviour. This statement is made in connection to the fact that there shall be creation or destruction of particles or complete obliteration of the system (blackhole evaporation) or obfuscation of results. Some systems are like "inside traders" they will not put signature atoll! How do you find they did it! Anyway, there are possibilities of a CIA finding out as they recently did! So we can do the same thing witn systems to. This is accentuation, corroboration, fortification, .fomentatory note to explain the various coefficients we have used in the model as also the dissipations called door. In the bank example we have clarified that various systems are individually conservative, and their conservativeness extends holisticallytoo.that one law is universal does not mean

there is complete adjudication of nonexistence of totality or global or holistic figure. Total always exists and "individual" systems always exists, if we donot bring Kant in to picture! For the time being let us not! Equations would become more ensorcelled and frenzied..... philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". In consciousness research, two rival sets of theories can be recognized: (A) Scientific material interpretations of consciousness are based on axioms that view consciousness in the context of highly advanced intentional processing of information in which subject-object relations evolve, and (B) humanistic interpretations of consciousness are based on axioms that view consciousness in the context of, say, "centered pulsations" that enable a conscious agent to act from his or her center of awareness. In this paper I will argue for the selection of axioms that favor humanistic interpretations of consciousness. The Pursuit of Autonomy Interdisciplinary Observations to Human Consciousness: Wautischer, Helmut (1993) the Pursuit of Autonomy. Interdisciplinary Observations to Human Consciousness: Social Neuroscience Bulletin, 6 (4). pp. 52-56.Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality; **CONSCIOUSNESS AND ITS UNIQUENESS June 12, 2012 at 8:40pm (See Deleuze Logic of sense, Wikipedia and Stanford encyclopedia for more details)** Atrocious contrivance and device, gratificational primogeniture, calamitous dodge and expedient, depraved, destructive, disastrous, execrable gambit and gimmick, great solace and succor, iniquitous, injurious, loathsome, low, maleficent, malevolent machination and maneuver, promethaleon of candor, frankness, honesty, honor, ingenuousness, innocence, openness, reality, sincerity, truthfulness, spiteful, stinking, ugly, unpleasant, unpropitious play and ploy, progenitor of jurisprudence and circumspection, racket and ruse, savvy, scam, stratagem, subterfuge, tactic and wile, forthrightness, honesty, truthfulness embodied and personified is the structural predisposition and dispensation of the complex. Nobody need have to take the Brahman-Anti Brahman agency to be a recalcitrant, repugnant, and refractory proposition. God gives and also takes away. That is the point made. ‘

And of course lesser gods, and God doth follow them and they follow God. That is why they are lesser gods. And wield the power on lesser mortals like us. When we mean gratification and deprivation we mean celestial Brahman Anti Brahman and Terrestial Brahman and Anti Brahman agencies producing such happiness or sadness by any means such as (blending- amalgamation, adulterant-adulteration, amalgam-amalgamation, blend- combination,

composite- compound, debasement-denaturant, fusion-hybrid, intermixture-reduction, amalgam- mixture, composite- compound, fusion-mishmash, aggregate- alloy, amalgamation-blend, commixture-composite, composition-compost, conglomerate-fusion, goulash-medley, mishmash-stew, blend-coalition, commixture-compound, federation-heating, integration-junction, melting-merger, synthesis-unification, assortment- combination, adulteration-alloy, assimilation-association, batter-blend, brewed combine, composite compound, concoction confection, conglomeration cross of money, orientation(dis), penury creation, murder, mayhem plunder, pillage, apocalypse, Armageddon in various forms and permutation and combination of Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness: with each action accounted and showing balance in the nature's general ledger as also actionlessness:e). All these are recorded in Nature's general ledger. Neuron DNA encapsulates all the actions passions, actions, interactions, transactions or the lack of it at various levels of individual, collective and cosmic levels. In celestial Brahman Anti Brahman ledger there is cessation of dualities. One more point to be made is if 5=3+2, then we can say 5 gormandizes 3 or 5 gobbles up 2. Quintessentially such a statement arises out of the difference of LHS and RHS term wise. It must be remembered despite all theoretical abstractions and generalizations be it proton mass or Higgs Boson, or topology of a space, it is notwithstanding the fundamentality of nature is just a number. Just everything in Physics and mathematics is just a number. Let us assume I buy rice in Bangalore (District) worth 5k.g. I can always divide 5k.g.s of rice in categories of Bangalore

(South), Bangalore (west), Bangalore (East) and Bangalore (South). On similar basis, there are billions of neutrons, protons and concomitant interactions which would satisfy the given numerical value. Example of gravity above makes things clear. Everything atleast is a function of time and someplace or some other coordinates, subordinates or superordinates. Mass of proton is 2.47 does not mean total mass of all protons is 2.47. It is this argument that is used in the classification scheme. By the very nature, self the unmonitoring but the witness consciousness agency identifies itself with something, be it a profession, an actor on the screen, when it does not have a identity. Such examples are legion, a Charlie Chaplin in "The Kid" or Raj kapoor in "Jagte Raho". These are people who just wanted to live and living without being put under surveillance every moment is the greatest part of their life. This is "Absolute Subjectivity" of Shiva. Such a state is always craving for identification for in pure consciousness all identities are lost. It is this "identitylessness" or having very minimal identity is what pervades most in India and this is due to expansion of individual consciousness. When that happens, be it education, or wealth, the "identity' becomes so well entrenched, they become the exact opposite of "pure consciousness"," impure consciousness" (Shakti). This is a state of "Absolute or relativistic objectivity". There was this case of a lady who went to such an extent of planning, she was sure that whatever she does must work out. When it did not, she lapsed in to "identity crisis". Such thing as this might happen in well developed countries where the competition is extremal and to be achieved at any cost. Realisation that one is not the doer but the witness consciousness is just like watching a movie in which one has acted. It is in this sense "advaita", "non duality' seems to be of cardinal and preeminent importance for the analytical treat ment of "identity' and "identity lessness", both happen in space time. "Multiverses" in this sense is also like a different film of an actor released at different theaters and it is in this sense we state that that every object in space time is a multiverse. You must have totally immersed in doing a "string theory" problem or watching a "film", oblivious of surroundings. Essentially creation (space time) also is what is happening in mind and honed in to you billion times. Identification of anything with a designation or signification is one that leads to

association within spacetime, which would only have pathological consequences and detrimental ramification as stated above or positive implications of "happiness" because of the identity itself. Search or 'show" has to go on! It gives a sense of consummation. a watching pot never boils; donot go to check nature; you are disturbing it; it will never show its true colours bhakti and Virakthi, gratification and deprivation, raga and dwesha is essential to live in space. It is the quantum interest thereof that produces energy to live .gratification and deprivation which form the bastion, pillar, post and stylobate and sentinel of all human actions depends on tams, rajas, and sattva of individual or insentient object. Verily tamas, rajas and satva are the characterstics of space itself, nay the behavioural pattern and attitudinal orientation. There are peaceful places like Vienna or Sringeri and there are violent places like others. Space is warped and woofed by the Akshara, the quantum information of primordial syllable Om. cosmic general ledger is the one vibration of which is brahman; quantum information possessions thereof form the sentinel of all transactions (Gargi to Yajnavalkya: my interpretation)adhocism, nonseqientiality of events, anti sequentiality of events, simultaneity of phenomenon, concurrent happenings , improbability of phenomenon happening itself are characterstics of brahman-anti brahman (Terrestial and celestial) are the cause of non predictability of happenings like in black hole, or in the areas of operation of brahman-anti brahman (Terrestial and celestial). brahman AntiBrahman actions, interactions, decisions, circumspection, adhocism, satisfies sensitive to initial conditions; topologically mixing; dense periodic orbits, the axiomatic predications and postulation alcovishness of chaos theory. Like gravity is a characteristic of space, so is tamas, rajas, sattva, with spaces knotted together by quantum information (for models see below) .accidents are preplanned actions of Terrestial and celestial Brahman AntiBrahman. All actions are perpetrated by them, be it train on fire or death. Belief systems play an important role in the determination of tamas, rajas, and satva state of human beings. There is neither cause nor effect. All actions are performed by Terrestial and celestial Brahman anti Brahman and it is off the record. psychic energy is negative equivalent to –mc^2, so that second law of thermodynamics is satisfied on a holistic basis vis-à-vis material energy

ParaBrahaman is infinite individual, collective and cosmic consciousnesses are infinite (concomitant general ledgers) for they are in the evolution mode with respect to each other in that order. Later emanates from the former as transactions emanate from the conservative state of assets and liabilities. Is not zero true for infinite conservative systems inside space and time? zero is the only consummation, consolidation and conservation. if fifty out of seventy members come out of the room at whatever time you come out, talk about the grand limbo of oblivion and hibernation of lifeless world you are going to and if you are not a celebrity, then it is not coincidence, it is coordination. Nature is playing the same trick on you. That a generalised law is there and you can measure does not make things less good albeit recalcitrant, refractory and rogue particles occur time and again, by and large. Entering in to all living beings through his constant part on the new moon night, he is born therefrom next morning. Absolute ParaBrahaman (zero) projects himself again and again on the screen of individual consciousness which appears of depairing despondent world. Some where the balancing act is done on karma. cosmic general ledger reflects upon itself to keep all transactional ties conservative (Brihadarnkayaka Upanishad: 1.5 14) for models see next paper through its divine power ( i interpret man creates god; it is attribution and ascription of various Shaktis would the formation of god takes place like space and time) the self assumes (eb) the form of a deity man can contemplate and venerate, even though siva, the pure subject, can never in fact be an object of meditation (adhyeya). Until we realise our true identity with Sankara, he is worshipped and conceived to be a reality alien to ourselves. While we are in the realm of creation, he too is a creation or mode or appearing of the absolute, manifest to us in meditation, through his freedom as an eternal, omniscient being. There is no gulf between the created and the uncreated creator: nothing in reality, although an object of knowledge, ceases to be Siva: this is the reason why meditation [on this or that aspect] of reality bestows its fruit. The world of the senses and mind appears to the well awakened (suprabuddha) as a theophany, an eternal revealing of god in his creation. The doctrine of vibration declares that "there is no state in word, meaning or thought, either at the beginning, middle or end, which is not Siva." 59 to utter any word is, in

reality, to intone a sacred formula.every act is a part of Siva's eternal cosmic liturgy, every movement of the body a ritual gesture (mudras), and every thought, god's thought. By what path are you not attainable? What words do not speak of you? in which meditation are you not an object of contemplation? What indeed are you not, o lord? 61 spandashakti, which accounts for the appearing of all things, is also the means by which deity in its many varied forms appears to man. Ksemaraja concludes: the ultimate object of worship of any theistic school differs not from the Spanda principle. The diversity of meditation is due solely to the absolute freedom of Spanda. 62 Sankara is not only the supreme object of devotion; as the static polarity of the absolute, he is the inner reality which holds together its Siva and Sakti outer manifestations. Phenomena are patterns of cognitions projected onto the surface of self-luminous siva-consciousness. There they become apparent, directly revealed to consciousness according to their manifest form. Siva is accordingly symbolized as the ground or surface of awareness, smooth and even like a screen (samabhittitalopama). 64 inscribed on this screen (kutfya) are the countless manifest forms which appear within it rendering it as diverse and beautiful as a fossil ammonite (falagrama). 65 Siva is the sacred ground upon which the cosmic mandalas are drawn, the absolute surface of inscription which bears the mark (cihna) of the universe. Abhinava writes: the variety of this world can only be manifest if the highest lord, who is essentially the pure light of consciousness, exists; just as a surface is necessary for a picture. if external objects were perceived in isolation then, because 'blue' and 'yellow', etc., are self-confined and the perceptions [we have of them] refer to their objects alone and so are insentient, mute and dumb in relation to one another . . . how would it be possible to be aware that an object is variegated? but just as depths and elevations can be represented by lines on a smooth wall, and we perceive [a female figure and think], similarly it is possible to be aware of differences in the variegated (contents of experience) only if all the diverse perceptions are connected together on the one wall of the universal light of consciousness. 66 Siva is the perfect artist who, without need of canvas or brush, paints the world pictures. The instant he imagines it, it appears spontaneously, perfect in every respect. The colours he uses are the varying shades

and gradations of his own Spanda energy and the medium his own consciousness. Intractable disaffect, Mutinous obdurate, Rebellious recalcitrant, Seditious refractory, Insubordinate insurgent, antagonistically antipodean, discrepant incompatible, Contumacious coadjutant, Conflicting converse, Unyielding obdurate, the antagonistic and anachronistic dispositional ties of Anti Brahman are in fact the sine qua non of Brahman only (See Vishwaroopa of Mahavishnu, which contains the polarised positives on RHS and negatives on LHS). The universe is coloured with the dye of its own nature (svabhava) by the power of Siva's consciousness (a//). Rajanaka rama says: homage to him who paints the picture of the three worlds, thereby displaying in full evidence his amazing genius {pratibha)\ to Sambhu who is beautiful with the hundreds of appearances laid out by the brush of his own unique, subtle and pure energy. 68 analogously, at the microcosmic level, all the cognitions and emotions, etc., which make up the individual personality form the outward flow of essentially introverted consciousness. They are specific pulsations (vise$aspanda) or aspects of the universal pulsation (samanyaspanda) of pure t consciousness. at the lower level, within the domain of maya, they represent the play in the fettered soul of the three primary qualities (gunas) or feeling-tones' which permeate to varying degrees his daily experience. These are: 1) sattva — the quality of goodness and luminosity which accompanies blissful experience both aesthetic and spiritual. 2) Rajas — the passion or agitation which oscillates between the extremes of 'light' and 'darkness' and characterises inherently painful experiences. 3) Tamas — the torpor and delusion which accompany states of inertia and ignorance. 69 the liberated soul recognises that these three are the natural and uncreated powers of pure consciousness. for him they are manifest respectively as: 1) Sankara 's power of knowledge (jnana) — the light of consciousness (prakasa); 2) the power of action (kriya) — the reflective awareness of consciousness (vimarsa)\ 3) the power of maya — which does not mean here the world of diversity, but the initial subtle distinction which appears between subject and object in pure consciousness. 70(doctrine of vibration: mark S.G. Dyczkowski) .to "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new

concepts and meanings that could not be included in the old system. at the heart of western metaphysics, for example, derrida finds the opposition between "speech" and "writing." this binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior in derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. The postmodern turn in philosophy: theoretical provocations and normative deficits by Steven best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html. You cannot see beyond consciousness for that is the limit of your storage of information or the basis for understanding further information. Turiya is empty bucket; it is from which prajna the measure of  arises; turiya measures the number of quantum information stores in other three states of waking, dreams, and dreamless deep sleep (Manduka Upanishad; my interpretation for models next paper).let there be path where angels tread: it is a force never born never dies the brahman an anti brahman verily there are parallel universes in (eb) in the ambit of space and time; the society of sinners and the sinned with the grammar of evil; parallel universes contain(e) multiverses in them; all physics, philosophy,  break down(e) in these universes; astronomical gratification ,deprivation, glorification, mortification, projections, axes, dimensions, rotations, foldings, concurrences, motions, shocks, are the axiomatic predications of multiverses and parallel universes; death is when self (the witness consciousness; register of individual general ledger) leaves the body; ostensibly parallel seem to have (e) an objective; parallel universes donot (e) have objectives; their paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common

attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, affection- appreciation, awareness- discernment, emotion-feeling, gut reaction, heart insight, intuition, judgment, keenness, perceptiveness, rationale, sensation, sense, sensitiveness, sensitivity, sentiment, susceptibility, taste, vibes, admiration, aesthetic sense, affection, appraisal, assessment, attraction, awareness, cognizance, commendation, comprehension, enjoyment, esteem, estimation, grasp, high regard, knowledge, liking, love, perceptional realization, recognitional regard, relish, respect, responsiveness, sensibility, sensitiveness, sensitivity, sympathy, understanding, valuation, acknowledgment, airing, baring, , confessional defenselessness, denudation- denunciation, disclosure, display, divulgence-exhibition, hazard, introduction, jeopardy, laying open, liability, manifestation, nakedness, openness, peril, presentation, publicity, revelation, risk, showing, susceptibility , susceptiveness,susceptivity, unfolding, unmasking, unveiling, vulnerability, vulnerableness, action, affection, appreciation and ardor, behavior, capacity, compassion, concern, cultivation and culture, delicacy, discernment and discrimination, emotion, empathy, faculty, fervor, fondness, heat, imagination, impression and intelligence, intensity and intuition, judgment, keenness- palpability, passion, pathos, pity, reaction, refinement, sensibility, sensitivity, sentiment, sentimentality , sharpness, spirit, sympathy, tangibility, taste, tenderness, understanding, warmth, fond memories, hearts and flowers, homesickness, longing, pining, reminiscence, ,schmaltz, sentimentality wistfulness, yearning satiation is only ostensible; their job is to get the job done; their business is to mind other's business ;all laws of physics and finance breakdown; it is antilaw,piscatorial Piratish and bubonic bucaneerishness; there is free transactions between in multiverses; when reflected on the screen of consciousness they appear as extant existential universe. Transactions are astronomical and gratification

deprivation, glorification and mortification incalculable; they have no barriers of region, religion, language which they create for others. theirs objectives are: corruption criminality, debasement debauchery, degradation evil immoral lewdness licentiousness, perversion sinfulness, vice viciousness, vileness wickedness, improbity corruption, atrocity decadence, degeneration degradation, depravity evil, immorality impurity, infamy iniquity, looseness lubricity, perversion profligacy, sinfulness turpitude,vice viciousness, vulgarity wickedness,disgrace bad reputation, abasement abuse, baseness black eye, blemish blur, brand comedown,contempt contumely, corruption culpability, debasement debasing, defamation degradation, derision disbarment, discredit disesteem, disfavor dishonor, disrepute disrespect, humbling humiliationignominy ill repute, infamy ingloriousness, meanness obloquy, odium opprobrium, pollution ,put-down reproach, scandal scorn, slander slight, slur spot, stain stigma, taint tarnish, turpitude venality, enormity Horribleness , abomination atrociousness, atrocity crime and the positive aspectionalities thereof .let there not be any incursions where evil mercenaries in the garb of mendicants thrive with adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatious gainsaying or conformingly pugnacious Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, execrable and goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness . statement for

Mephistophelean mercenaries philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality. Negative energy stabilising wormholes are considered candidates for (er) darkmatter. An analysis of negative energy might be a useful ingredient of the thoughts of travel at speeds more than the velocity of light. Peirce explains that an "imperfect continuum" possesses "topical singularities, or places of lower dimensionality" which break continuity, while "the parts of a perfect continuum have the same dimensionality

as the whole." (CP 4.642) Circular continua without singularities are time and space. (CP 6.210-212) By using math as his lens, Peirce envisioned singularities or lower dimensional holes in the imperfect cosmologic continuum, while maintaining that a probabilistic law of chance exists that entropy decreases **(Wikipedia).** On the side of philosophy**,** quite explicitly, Deleuze provides a critique of Husserl's "sleight-of-hand" change from a static categorical phenomenology to one supposedly more dynamically constitutive of reality. This a clear cut **chagrined circumvention**, **contraventional curbing**, **defeated disgruntlement**, **dissatisfactional downer**, **dragged failure**, **fizzled foil**, **grievance hindrance**, **impedimental irritation**, letdown, nonfulfillment, nonsuccess, obstruction, old one-two, **resent mental setback**, unfulfillment of the earlier theories makes Sartre contend that human existence is a conundrum whereby each of us exists, for as long as we live, within an overall condition of nothingness (nothing-ness)—that ultimately allows for free consciousness. But simultaneously, within our being (in the physical world), we are constrained to make continuous, conscious choices. It is this dichotomy that causes anguish, because choice (subjectivity) represents a limit on freedom within an otherwise unbridled range of thoughts. Subsequently, humans seek to flee our anguish through action-oriented constructs such as escapes, visualizations, or visions (such as dreams) designed to lead us toward some meaningful end, such as necessity, destiny, determinism (God), etc. Thus, in living our lives, we often become unconscious actors—Bourgeois, Feminist, Worker, Party Member, Frenchman, Canadian or American—each doing as we must to fulfill our chosen characters' destinies. However, Sartre contends our conscious choices (leading to often unconscious actions) run counter to our intellectual freedom. Yet we are bound to the conditioned and physical world—in which some form of action is always required. This leads to failed dreams of completion, as Sartre described them, because inevitably we are unable to bridge the void between the purity and spontaneity of thought and all-too constraining action; between the being and the nothingness that inherently coincide in our self. Deleuze insists on the absolute non-resemblance between what conditions (metastable, pre-individual plane of singularities) and what is conditioned

(actualized or the individuated 'thing'). Only then do "the conditions of the true genesis become apparent," Deleuze writes (Logic of Sense 105). Without addressing this issue, Deleuze's borrowings from Simondon, Lautman and structuralism remain untethered. Because of Bowden's decision to completely sidestep Deleuze's overt negotiations with phenomenology, his otherwise sophisticated commentary suggests a pedagogical narrowness. If, the omission of phenomenology strengthens Bowden's focus on the ontological priority of the event over substantial ontologies via the genealogy of figures and movements Deleuze outlines, it also incompletely articulates the problem of the sense-event as Deleuze formulates it. We have lost any idea of whom Deleuze takes to be his primary adversary, against whom he stakes out his position. This worry is most clearly manifested in Bowden's description of Deleuze's ontological prioritizing of events as offering a "transcendental ontology." For Bowden Deleuze's transcendental ontology defines events that "are ontologically prior to substances, 'all the way down'" (82). But Bowden provides a second definition: it is that the priority of events over worldly individuals means that the world of events "is not something external to the conditions of knowledge" (69). **SEAN BOWDEN the Priority of Events: Deleuze's Logic of Sense Sean Bowden, the Priority of Events: Deleuze's Logic of Sense, Edinburgh University Press, 2011, 296pp., $40.00 (hbk), ISBN 9780748643646. Reviewed by David Scott, Coppin State University Weltschmerz, agony, apprehension, blues, depression, dread, mid-life crisis, misgiving, nervousness, uneasiness, angst were expressed with such effulgent words…… aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement, argument, articulational assertion, asseverational choice of words, commentary and communication, declaration and definition, delivery, diction, elucidation, emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation about the human predicament is brought out in Being and Nothingness. Being and Nothingness offers a critique of Sigmund Freud's theory of the unconscious, based on the claim that consciousness is essentially self-conscious. Sartre also argues that**

**Freud's theory of repression is internally flawed. According to Sartre, in his clinical work, Freud encountered patients who seemed to embody a particular kind of paradox—they appeared to both know and not know the same thing. In response, Freud postulated the existence of the unconscious, which contains the "truth" of the traumas underlying the patients' behavior. This "truth" is actively repressed, which is made evident by the patients' resistance to its revelation during analysis. Yet what does the resisting if the patients are unaware of what they are repressing? Sartre finds the answer in what Freud calls the "censor". "The only level on which we can locate the refusal of the subject," Sartre writes, "is that of the censor.** The Holographic Principle, as conceived in current physics, applies to fields, and perhaps, even to more elementary entities, called strings and branes. These more elementary entities remain hypothetical at the time of this writing. There are many layers between the level of fields and that of neurons. We know that neurons are surrounded by a surface phospholipid membrane, which supports an electrochemical process that is fundamentally necessary to human experience. As the information supported by the membrane surface according to the Holographic Principle is proportional to its area, we should expect to find that the neural correlates of consciousness most clearly associated with those parts of the neuron that have the have a high proportion of neuronal surface area, and, in particular, receptive surface area. Dendrites account for an average of about 90% of the neuron is receptive surface area (Wong, 2002), and so would be expected to be the most important neural structures in information processes. We find this to be true **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science (Wikipedia).** Discordant-discrepant, dissonant-hard-line, implacable, incompatible, incongruous- inconsistent, inexorable-inflexible, inharmonious-intransigent, opposed, reluctant, unappeasable, uncompromising, unfriendly some of these expatiations and enumerations might sound, a more generalised interpretations is given by a Vedantin. Look at this arresting passage: The teacher continued to be silent. When addressed a second and third time he said: "I am teaching, but you do not follow. The Self is silence." The undetermined and unthinkable character of the Brahman is

a consequence of the absolute's eternal and immutable nature. To concede the existence of a real universe is, from the Vedantin's point of view, to posit the existence of a reality apart from the Brahman. Nor can we simply identify a real universe with the absolute unless we are prepared to compromise its unchanging, absolute status. The criterion of authenticity is immutability. Reality never changes; only that which is less than real can appear to do so. Reality is constant in the midst of change. What this means essentially is that there is change although nothing changes. This impossible situation is reflected in the ultimate impossibility of change itself. That which does not exist prior to its changing and at the end, after it has changed, must be equally non-existent between these two moments. Although the world of change appears to be real, it cannot be so. Change, according to the Vedantin, presupposes a loss of identity. Reality cannot suffer transformation; if it were to do so, it would become something else and the real would be deprived of its reality. The immortal can never become mortal, nor can the mortal become immortal. The ultimate nature of anything cannot change. Change of any sort is merely apparent (vivarta); the world of change and becoming is a false super- imposition (adhyaropa, adhyasa) on the absolute. In cosmic terms, the mistake (bhranti) consists of the supposition that the real Brahman is the unreal universe and the unreal universe is the real Brahman. In microcosmic terms, it is the mistake of falsely conceiving the body, mind or even one's personality to be the Self. In the same way as the image of a snake is falsely superimposed on a rope, similarly the universe is falsely projected onto the real substratum, the Brahman. Ignorance is not merely a personal lack of knowledge, but a cosmic principle. As such it is called "Maya," the indefinable factor (anirvacaniya) that brings this mistake in identity about. The reality status of this cosmic illusion is also indefinable: on the one hand it is not Brahman, the sole reality; on the other hand it is not absolutely non-existent like a hare's horn or the son of a barren woman. Brahman is the source of world appearances only in the sense of being their unconditioned ground or essential nature. The universe is false not because it has no nature of its own but because it does have one. Just as the illusion of a snake disappears when one sees that it is nothing but a rope, similarly cancellation (badha) of the empirically real occurs

when the absolute reality of the Brahman is realised. Thus, according to Vedanta, appearance implies the real, while the real need not imply appearance. To appear is essentially to appear in place of the real, but to be real is not necessarily to appear. All things -exist because the absolute exists. It is their Being. Thus the very existence of phenomena implies their non-existence as independent realities.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^\infty \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9 \left( U_9 - V_9 \int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9 \right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9 \int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$then given the constraints of $f(T_9), (A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 -$$

$$V_9 M_9 -]e^{-T_9\lambda_9} = 0$$

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9 Z_9 \sin(Z_9 T_9) - (A_9 M_9 + Y_9) + (A_9 M_9 + X_9)\cos(Z_9 T_9) = 0$$

$$(A_9 M_9) Z_9 - M_9 Z_9 \cos(Z_9 T_9) + (A_9 M_9 - X_9)\sin(Z_9 T_9) = 0$$

Quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High

Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) **Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP** This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) **Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP** In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference: Journal of Physics: Conference Series 70 **Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows:** System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol

Where $A_9 = U_9 - V_9$Using (A2.3), substitute for $\boldsymbol{Z}_9^2$ and $Z_9$in (A2.2):

$W_9^2 (X_9^2 - 3 M_9 A_9 X_9 - M_9^2 (X_9 - Y_9) + 2 M_9^3 A_9 + 2 A_9^2 M_9^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_9$=1-cos ($Z_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (X_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ X_9 - Y_9/2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \leq W_9 \leq 2$ *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (X_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\,M_9)^2\,(2A_9 M_9 - (X_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9 M_9^2 - 3M_9 X_9\right) + W_9\left(3M_9 A_9^2 - 2A_9 X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9+Y_9\right)$$

$$=2M_9^2W_9^2(M_9+\mathrm{A}_9)+M_9A_9^2W_9+\left(2A_9M_9-X_9+Y_9\right)(A_9W_9+M_9W_9^2$$

$$+(2-W_9)((2M_9X_9+M_9Y_9)W_9+A_9(X_9+Y_9)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range 0 $W_9$ 2.Defining

$Q_9\left(W_9,A_9\right)=a_9W_9^2+c_9$ *where* $a_9b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

To Heidegger writings we shall compose the model. Following is not still written model able form. For one of the models please see any one of the papers under the appellation: Ïn Excited Reveries

In his later writings Heidegger drops the term 'phenomenon' in preference for the verbal form 'phainesthai' in order to emphasize even more the actuality or presentational property of Being. Explaining this new form of the term he writes: "Being disclosed itself to the ancient Greeks as 'physis'. The etymological roots 'phy-' and 'pha-' designate the same thing: 'phyein', the rising-up or upsurge which resides within itself as 'phainesthai', lighting-up, self-showing, coming-out, appearing-forth." Heidegger contrasted his notion of phenomenon with semblance (Schein) and with appearing (Erscheinung). In the case of semblance a thing can show itself as that which it is not, as when fool's gold shows itself to be gold. The ancients always allied semblance with non-being. Heidegger points out, however, that semblances are grounded in

showings, and so does Abhinava. Both Heidegger and Abhinava consequently maintain that all semblances have a real basis and are to be treated as instances of phenomena along with the so-called real showing or manifestation of non-deceptive objects. So Heidegger states that: 'how- ever much seeming, just that much being'. Thus self-showing or appearing defines Being as phenomenon, but this definition of Being is as yet incomplete. Being is not only self-showing but 'logos' which Heidegger explains means 'discourse' (Rede) in the sense of 'apophansis': 'letting-be-seen'. Phenomenology, which according to Heidegger is the only correct study of Being, means 'letting-be-seen-that-which- shows-itself. This is true of Saiva Paramadvaita as well. The reality of the world demands recognition; we are forced to accept the direct presentation of the fact of our daily experience. As Abhinava says: "if practical life, which is useful to all persons at all times, places and conditions were not real, then there would be nothing left which could be said to be real" A thousand proofs could not make 'blue' other than the colour blue. The reality of whatever appears in consciousness cannot be denied. Objects appear; they do not cease to do so by a mere emphatic denial. The manifestation of an entity in its own specific form is a fact at one level of consciousness; it is real. The appearing of the same entity in the same form but recognised to be a direct representation of the absolute is also a fact, but at another level of consciousness. It is no more or less real than the first. 'As is the state of consciousness, so is the experience,' says Abhinava. Although the nature of the absolute is discovered at a higher level of consciousness, The Doctrine of Vibration nonetheless it presents itself to us directly in the specific form in which we perceive things; otherwise there would be no way in which we could penetrate from the level of appearing to that of its source and basis. Abhinava writes: Real is the entity iyastu) that appears in the moment of direct perception (sak$atkara), that is to say, within our experience of it. Once its own specific form has been clearly determined one should, with effort, induce it to penetrate into its pure conscious nature. All things are known to be just as they present themselves. The concrete actuality of

being known (pramiti), irrespective of content, is itself the vibrant (spanda) actuality of the absolute. Liberating knowledge is gained not by going beyond appearances but by attending closely to them. "The secret," Mahesvarananda says, "is that liberation while alive (jivanmukti) is the profound contemplation of Maya's nature." No ontological distinction can be drawn between the absolute and its manifestations because both are an appearing (dbhasa), the latter of diversity and the former of 'the true light of consciousness which is beyond Maya and is the category Siva'. Those who have attained the category of Pure Knowledge above Maya and have thus gone beyond the category of Maya, see the entire universe as the light of consciousness . . . Just as the markings [on a feather] are nothing apart from the feather, the feather [is nothing apart from] them, similarly, when the light of consciousness is manifest, the whole group of phenomena is manifest as the light of consciousness itself. Within the sphere of Maya, every entity's 'own nature' (svabhdva) corresponds to its specific manifest form. Accordingly it is defined as that which distinguishes it from all else and from which it never deviates. Above the sphere of Maya, that is, above the level of objectivity, is the domain of the subject. At this level, everything is realised to be part of the fullness of the experience r and hence no longer bound by the conditions which impinge on the object. Here the part is discovered to be the whole, that is, consciousness in toto. In this sphere beyond relative distinctions, the yogi realises that (all) the categories of existence are present in every single category. m The yogi experiences every individual particular as the sum total of everything else. He recognises that all things have one nature and that every particular is all things. This is the 'essence' (sard) or co-extensive unity (samarasya) of all things. Integral Monism of Kashmiri SaivismWe have established that reality is manifest according to how [and the degree in which] the freedom of consciousness reveals it and that [this freedom] is the womb of all forms. Just as 'sweetness' is present in its entirety in every atom of the sugarcane, so each and every atom [of the universe] bears within itself the emanation of all things. This is

the level of consciousness in which the absolute reflects on itself realising to its eternal delight and astonishment (camatkara) its own integral nature. The reality of the world of diversity is not denied, but experienced in a new mode of awareness free of time and space in the eternal omnipresence of the Here and Now. [Phenomenal forms of awareness] such as 'this [exists]', born of the colouring [imparted to the absolute] by the limitations engendered by the diversifying power of time {kalakalana) also emanate within the Supreme Principle. There [at that level], Fullness {purriata) is the one nature [of all things] and so everything is omnipresent; otherwise, associated with division (khancjlana), the Fullness [of the absolute] would not be full. The content of absolute consciousness consists of diverse appearings (abhdsa) which, because they are manifest through it in this way, do not compromise the wholeness of consciousness. Everything we perceive is a momentary collocation of a number of such manifestations which combine together like 4 a row of altar lamps' (dipavalT) to form the single radiant picture of the universe. The individual objects which constitute the universe are specific collocations of such 'atomic' appearings. Together they form a single unified particular which appears according to its own defining features (svalaksand). A jar, for example, consists of a number of appearances such as 'round', 4 fat\ 'earthen', 'red', etc., which together discharge a single function (arthakriya), in this case, that of carrying the appearance 'water'. They unite with each other much as the scattered rays of a lamp come together when focused, or as the various currents of the sea together give rise to waves. Atomic appearings can combine in any number of ways, provided that they are not contrary to one another as established by the dictates of natural law (niyati). An appearance of 'form', for example, cannot combine with that of 'air'. Insofar as they share a common basis (samanyadhikaranya), a given cluster of appearances appears as a single whole. This common basis is the most prominent member of the group; the appearance 'jar' is such in the example quoted above. Any one appearance in a cluster may assume a more important or subordinate role.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_9 \square - b_9 \square 0$$

And

$$b_9^2 \square 4a_9 c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9 in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\ X_9$ ,

$Let\ Y_9 = \ y_9\ X_9$And try the substitution $A_9 = X_9/(2 + A_9\ y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + \ n_9 y_9)$ we have

$$W_9^2\ y_9\left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^2\left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9\,X_9\left[\frac{X_9\left(1+n_9y_9\right)}{M_9^2\left(2+n_9y_9\right)^2}+2\,y_9\left(n_9-1\right)+\frac{X_9^2\left(9+y_9\right)}{M_9^2\left(2+n_9y_9\right)}\right]=0.\ (A2.12)$$

Since $n_9$ and $A_9$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9=\left(2+n_9y_9\right),$ *the condition*$(A2.99?)$*can be written as*

$$\left[\frac{X_9\left(E_9-9\right)}{E_9^2}\right]^2\left[\frac{X_9^2}{M_9^2}-4\,y_9\,E_9^2\frac{X_9}{M_9^2}+4y_9\left(n_9-2\right)E_9^3\right]\ \square 0.$$

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9/W_9^2)$ can be expressed as

$$2\,E_9\left[y_9E_9+\left(y_9\,E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}\right]>\frac{X_9}{M_9^2}>2E_9\left[y_9E_9-y_9E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}]$$

dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique

for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) **Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP** Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) **Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP** Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a

cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) **Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP** Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service

placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) **Dynamic Service Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP** Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by $H*(X,(\wedge*TX)\otimes g)$, where X is the target hyperk\"ahler manifold. We further give a mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) **A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van**

**Raamsdonk1 and Nathan Seiberg2** study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. **Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP** Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy **The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP** Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the

present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. **ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP** Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition. Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb) the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537

© 2016 American Physical Society **Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP** Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant ($\Delta\alpha/\alpha$) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and $\Delta\alpha/\alpha$ to (e&eb) this conclusion. However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts $z\lesssim 3$, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the $\Lambda$ cold dark matter model value ($j\Lambda CDM=1$); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in $\alpha(z)$, linked to (e&eb) a linear coupling ($\zeta$) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order $|\zeta|\sim 10-4$, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in $\Delta\alpha/\alpha$ (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society **The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP** Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation

rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void

Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems, authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society **Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP** Authors measure the star formation properties of (e) two large samples of

galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the **star formation** rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a **function of environment**. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]**.** Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size $\sqrt{n} \times \sqrt{n}$? The problem is that $\sqrt{n}$ time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time $O(\sqrt{n})$ for $d \geq 3$, or $O(\sqrt{n}\log 5/2n)$ for $d = 2$. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search

tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ∗This work was mostly done while the author was a **PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship and by CIAR**. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 **Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP** Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can

lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 **Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP** Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. **Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP** Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e)

sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) **Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP** The relationship between the foundations of mathematics and (e&eb) physics is a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb) flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling

field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) **Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP** The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar

and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph] (or arXiv: 1503.05600v1 [math-ph] for this version) **Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP** This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs

mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) **Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP** Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and

derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to **the fine structure constant**. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) **Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP** Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor goes to 0 as the (e&eb) time goes to 0, the

big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph] (or arXiv: 1306.4613v1 [quant-ph] for this version) **Effects of gauge theory based number scaling on geometry Paul Benioff STOP** In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x,

outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph]

(or arXiv: 1211.3381v2 [quant-ph] for this version) **Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP** Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value,

\int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances, the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) **Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP** The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include

(e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) **Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP** Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) **New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP** This work is based on (e) the field of reference frames based on (e) quantum

representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects:Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc) Journal reference:Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) **A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul B**

**ABSOLUTE STABILITY**

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 ? $W_9$ ? 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9-9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2-2y_9\frac{(E_9+n_9^2y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right)-4y_9(E_9-n_9)$$ (A2.14)

And has the solution

$$\frac{X_9}{M_9^2}>\frac{E_9}{(E_9-9)}[y_9(E_9+n_9^2y_9)+[y_9^2(E_9+n_9^2y_9)^2+4y_9(E_9-9)(E_9-n_9)]^{9/2}]$$

Given $E_9>n_9$this lower limit to $X_9/M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9/M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9<n_9$ $(ie.,2A_9M_9^2<X_9-Y_9)$then (A2.9)and (A2.8)show that $Q_9$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2}>2E_9[y_9E_9-(y_9E_9(9+y_9)(E_9-n_9)]^{9/2}]$$

And

$$E_9>n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2/2M_9(X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$ and $X_9$ and $Y_9$ by $W_9^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

**Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability

of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e) long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007 STOP. David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer P** examine the possibility of coherent reversible information transfer between solid-state superconducting qubits and (e&eb) ensembles of ultracold atoms. Strong coupling between these systems is mediated by (e&eb) a microwave transmission line resonator that interacts near resonantly with (e&eb) the atoms via (e&eb) their optically excited Rydberg states. Solid-state qubits can then be used to (e) implement rapid quantum logic gates, while (e) collective metastable states of (e) the atoms can be employed for (e) long-term storage and (e&eb) optical readout of quantum information. **Reversible state transfer between superconducting qubits and (e&eb) atomic ensembles David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer Phys. Rev. A 79, 040304(R) – Published 14 April 2009 STOP.** Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures have been proposed. An

attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. **Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier** demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in (eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. **Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP.** In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown

to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , **Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally** frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. **Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg** propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states

into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. **Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011 http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan** proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 **Qubit transfer between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller** propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. **Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys.**

**Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich** report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. **Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi** investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. **Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev** investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer.

Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. **L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP.** In recent years **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P have** investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P** apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys. Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. **Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S.**

**Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP.** Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert **Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe** realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.**220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2** [quant-ph] for this version) **Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP.** Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. **Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker** describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue

from over (e&eb) a thousand plant and animal samples. **NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP.** Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this behavior, enhancing and accelerating (eb, eb+) the growth of correlations. **Michael Kastner** gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. **Michael Kastner** tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) **Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP.** A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. **Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E. Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91,**

**075435 – Published 26 February 2015 STOP.** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still

problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to

(e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP** Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that

(eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035

**Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime**

**C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013**

**STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and

cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP** Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex

Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) **A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen**

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9 and$ separate (A3.1)into real and imaginary parts

$$Z_9^2 + \left(M_9 + V_9\right) Z_9 \sin\left(Z_9 T_9\right) - \left(U_9 M_9 + Y_9\right) + (M_9\left(U_9 + V_9\right) - X_9)$$

$$\mathrm{X}\cos\left(Z_9 T_9\right) - 2M_9 V_9 cos^2\left(Z_9 T_9\right) + M_9 V_9 = 0$$

$$(M_9 + U_9)\, Z_9 - (M_9 + V_9)\, Z_9\ cos\left(Z_9 T_9\right) + \left(M_9\left(U_9 + V_9\right) - X\right)$$

$$\mathrm{X}\sin\left(\left(Z_9 T_9\right) - M_9 V_9 \sin\left(Z_9 T_9\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 =$ $1 - cos\, Z_9\ T_9$and $A_9 = U_9 - V_9$we obtain

$$4\, M_9^2 V_9^2 W_9^4 + 2\, M_9 V_9\ W_9^3[(M_9 - V_9)^2 + A_9\left(M_9 + V_9\right) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2\left[M_9 - V_9\right)^2\left(2\, A_9\ M_9 + Y_9 - X_9\right) + \left(2A_9\ M_9 - X_9\right)^2 - A_9\left(M_9 - V_9\right)\left(2A_9\ M_9 - X_9\right)$$

$$+ M_9 V_9\,(A_9^2 + 4\left(X_9 + Y_9\right))]$$

$$+ W_9\left[A_9^2 + 2X_9\right)\left(2\, A_9\ M_9 - X_9\right) + 2A_9 Y_9\,(M_9 + V_9) - M_9 A_9^3] + A_9^2\left(X_9\ M_9 - Y_9\right) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9\left(2, U_9 V_9\right) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting cos $(Z_9 T_9) = -1\ in\ (A3.3)$

We have applied the model to Kashmiri Shaivism. Whose elaborate rituals is enough to make a indivual raise from individual consciousness to cosmic consciousness which we have called Cosmic Consciousness.

It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) **Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP** The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic

Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** The group of unimodular transformations on the roots of polynomials over GF(2) is considered, and those polynomials with (e&eb) symmetries in (eb) the unimodular group are identified. The cross-correlation function between (e&eb) two maximum-length linear shift register sequences of (e) the same degree is shown to be (=) computable as an explicit linear transformation, in (eb) matrix form, on (e&eb) either one of the sequences, regarded as (=) a vector. The underlying vector space is (=) the "cyclotomic algebra," generated by (e) the cyclotomic cosets, or by what Gauss termed the **"periods" of the cyclotomic equation**. A variety of numerical examples are worked in detail. **Information Sciences Volume 1, Issue 1, December 1968, Pages 87-109 Theory of transformation groups of polynomials over GF(2) with applications to linear shift register sequences Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more link. Solomon W. Golomb** stop Authors study the P−V criticality and phase transition in (eb) the extended phase space of charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of (e) the system and its conjugate quantity is (=) the thermodynamic volume of the black hole. The black holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb) no P−V criticality and phase transition appear, while (e) for the black holes with a spherical horizon, even when the charge of the black hole is absent, t(eb) he P−V criticality and the small black hole/large black hole phase transition will appear, but it happens only in d=5 dimensions; when the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to have (e) the P−V criticality

and (e&eb)the small black hole/large black hole phase transition, there exists (ebb) an upper bound for the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the **Gauss-Bonnet coefficient** and Q is (=) the charge of the black hole. Authors calculate the critical exponents at the critical point and find that (eb) for all cases, (eb) they are the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09(2013)005 Cite as: arXiv:1306.6233 [gr-qc] (or arXiv:1306.6233v4 [gr-qc] for this version) **P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in IR but with (e&eb) improved UV behaviors. Here, authors study an IR modification which breaks ``softly'' (e) the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Authors obtain the black hole and cosmological solutions for **``arbitrary'' cosmological constant** that represent (eb) the analogs of the standard Schwarzschild-(A)dS solutions which can be asymptotically (A)dS as well as flat and authors discuss their thermodynamical properties. They also obtain solutions for (e) FRW metric with (e&eb) an arbitrary cosmological constant. Authors study its implication to (e&eb) the dark energy and find that it seems to be consistent with (e&eb) current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP.** The Einstein equations with (e&eb) a negative cosmological constant admit (e) black hole solutions which are asymptotic to (e) anti-de Sitter space. Like black holes in asymptotically flat space, these solutions have (e) thermodynamic properties including (e) a characteristic temperature and intrinsic entropy equal to (=) one quarter of the area of the event horizon in Planck units. There are however some important differences from (e) the asymptotically flat case. A black hole in anti-de Sitter space has (e) a minimum temperature which occurs when (e) its size is of the order of the characteristic radius of (e) the anti-de Sitter space. For larger black holes the red-shifted

temperature measured at (eb) infinity is greater. This means that such black holes have (e) positive specific heat and can be in stable equilibrium with (e&eb) thermal radiation at a fixed temperature. It also implies (eb) that the canonical ensemble exists for asymptotically anti-de Sitter space, unlike the case for asymptotically flat space. One can also consider the microcanonical ensemble. One can avoid (e) the problem that arises in asymptotically flat space of (e) having to put the system in (eb) a box with unphysical perfectly reflecting walls because (e) the gravitational potential of anti-de Sitter space acts as a box of finite volume. Communicated by A. Jaffe **Communications in Mathematical Physics December 1983, Volume 87, Issue 4, pp 577-588 Thermodynamics of black holes in anti-de Sitter space S. W. Hawking, Don N. Page STOP** In previous work authors introduced the concept of black-hole entropy, which they identified with (e&eb) the surface area of the black hole in question expressed in units of (e) the Planck length squared. Authors suggested that the appropriate generalization of the second law for a region containing a black hole is (=) that the black-hole entropy plus (+) the common entropy in the black-hole exterior never decreases. Here authors establish the validity of this law for the infall of an entropy-bearing system into (e&eb) a much larger and more massive generic stationary black hole. To do this authors determine (eb) a general lower bound for the increase in black-hole entropy, and an upper bound for the entropy of the system, while allowing for quantum effects at each stage. In passing authors show that the generalized second law is a statistical law which becomes (=) over-whelmingly probable in the limit of a macroscopic system. Authors also consider briefly more general situations. Finally, we give two simple examples of predictions made by the generalized second law for (e) black-hole formation processes. Received 17 September 1973 DOI:http://dx.doi.org/10.1103/PhysRevD.9.3292 **Generalized second law of thermodynamics in black-hole physics Jacob D. Bekenstein Phys. Rev. D 9, 3292 – Published 15 June 1974 STOP** The Lovelock gravity consists of (e) the dimensionally extended Euler densities. The geometry and horizon structure of (e) black hole solutions could be quite complicated in (eb) this gravity; however, authors find that some thermodynamic quantities

of the black holes like the mass, Hawking temperature and entropy have (e) simple forms expressed in terms of (e) horizon radius. The case with black hole horizon being a Ricci flat hypersurface is (=) particularly simple. In that case the black holes are (=) always thermodynamically stable with (e&eb) a positive heat capacity and their entropy still obeys (e&eb) the area formula, which is no longer valid for (e) black holes with positive or negative constant curvature horizon hypersurface. In addition, for black holes in (eb) the gravity theory of Ricci scalar plus a 2n-dimensional Euler density with (e&eb) a positive coefficient, thermodynamically stable small black holes always exist in (eb) D=2n+1 dimensions, which are absent in (eb) the case without (e) the Euler density term, while (e) the thermodynamic properties of the black hole solutions with (e&eb) the Euler density term are qualitatively similar to those of black holes without (e) the Euler density term as D>2n+1. Copyright © 2004 Published by Elsevier B.V. **Physics Letters B Volume 582, Issues 3–4, 4 March 2004, Pages 237–242 A note on thermodynamics of black holes in Lovelock gravity Rong-Gen CAI STOP** There are a number of similarities between black-hole physics and (e, eb, =) thermodynamics. Most striking is the similarity in the behaviors of black-hole area and (e&eb) of entropy: Both quantities tend to increase (eb+) irreversibly. In this paper we make this similarity the basis of (e) a thermodynamic approach to black-hole physics. After a brief review of the elements of the theory of information, authors discuss black-hole physics from the point of view of information theory. Authors show that it is natural to introduce (eb) the concept of black-hole entropy as (=) the measure of information about (e&eb) a black-hole interior which is (=) inaccessible to an exterior observer. Considerations of simplicity and consistency, and dimensional arguments indicate (eb) that the black-hole entropy is equal to (=) the ratio of the black-hole area to the square of the Planck length times a dimensionless constant of (e) order unity. A different approach making use of the specific properties of Kerr black holes and of concepts from information theory leads to (eb) the same conclusion, and suggests (eb) a definite value for the constant. The physical content of (e) the concept of black-hole entropy derives from (e) the following generalized version of the second law: When common entropy goes down a

black hole, (eb) the common entropy in (eb) the black-hole exterior plus the black-hole entropy never decreases. The validity of this version of the second law is supported by (e) an argument from information theory as well as by several examples. Received 2 November 1972 DOI:http://dx.doi.org/10.1103/PhysRevD.7.2333 **Black Holes and Entropy Jacob D. Bekenstein Phys. Rev. D 7, 2333 – Published 15 April 1973 STOP** The effects of power-law Maxwell field on the Van der Waals like phase transition of higher-dimensional dilaton black holes are probed in detail. It is shown that the Smarr relation gains corrections due to (e) the effects of (e) both the power-law Maxwell field and (e&eb) the dilaton field while (e) thermodynamic volume is exactly the same as (=) that of Einstein-Maxwell-dilaton black holes. Authors successfully derive the analytic solutions of (e) critical point and carry out (e&eb) some check to ensure that these critical quantities are (=) positive. It is shown that the constraint on the parameters turns out to be (=) 0<\alpha^2<1, which is tighter than that in (eb) the non-extended phase space. It is also shown (eb) that these critical quantities and (e&eb) the ratio P_cv_c/T_c are affected by (e&eb) the power-law Maxwell field. Moreover, critical exponents are found to coincide with those of other AdS black holes, showing (eb) the powerful influence of mean field theory. Comments: references added Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1601.05500 [gr-qc] (or arXiv:1601.05500v2 [gr-qc] for this version **Effects of power-law Maxwell field on the Van der Waals like phase transition of higher dimensional dilaton black holes Jie-Xiong Mo, Gu-Qiang Li, Xiao-Bao Xu STOP** A consistent and unified picture for (e) critical phenomena of charged AdS black holes in f(R) gravity is drawn in this paper. Firstly, we investigate the phase transition in canonical ensemble. Authors derive the explicit solutions corresponding to (e&eb) the divergence of CQ. The two solutions merge into (e&eb) one when (e) the condition Qc=−13R0−−−√ is satisfied. The curve of specific heat for Q<Qc has two divergent points and can be divided into (e&eb) three regions. Both the large radius region and the small radius region are (=) thermodynamically stable with (e&eb) positive specific heat while (e) the medium radius region is unstable with (e&eb) negative specific heat. However, when Q>Qc, the specific heat is

always positive, implying (eb) the black holes are locally stable and no phase transition will take place. Secondly, both the T−r+ curve and T−S curve f(R) AdS black holes are investigated and they exhibit Van der Vaals like behavior as the P−v curve in the former research. Critical physical quantities are obtained and they are consistent with (=) those derived from the specific heat analysis. Authors carry out numerical check of Maxwell equal area law for the cases Q=0.2Qc, 0.4Qc, 0.6Qc, 0.8Qc. The relative errors are (=) amazingly small and can be negligible. So the Maxwell equal area law holds for (e) T−S curve of f(R) black holes. Thirdly, authors establish geometrothermodynamics for (e) f(R) AdS black hole to examine the phase structure. It is shown that the Legendre invariant scalar curvature R would diverge exactly where (e) the specific heat diverges. To summarize, the above three perspectives are consistent with (e&eb) each other, thus providing (eb) a unified picture which deepens the understanding of (e) critical phenomena of (e) f(R) AdS black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1602.01251 [gr-qc] (or arXiv:1602.01251v1 [gr-qc] for this version) **A consistent and unified picture for critical phenomena of f(R) AdS black holes Jie-Xiong Mo, Gu-Qiang Li, Yu-Cheng Wu STOP** The state equation of the charged AdS black hole is reviewed in (e&eb) the T−r plane. Thinking of the phase transition, the T−S, P−V, P−ν graphs are plotted and then the equal area law is used (e) in the three cases to get (eb) the phase transition point (P, T). The analytical phase transition point relations for (e) P-T of charged AdS black hole has been obtained successfully. By comparing the three results, authors find that the equal area law possibly cannot (e) be used directly for P−ν plane. According to the T−S, P−V results, they plot (e&eb) the P−T−Q graph and find that for a highly charged black hole a very low temperature condition is required for (e) the phase transition. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI:10.1140/epjc/s10052-015-3641-0 Cite as: arXiv: 1503.07658 [gr-qc] (or arXiv:1503.07658v1 [gr-qc] for this version) **A note on Maxwell's equal area law for black hole phase transitionShan-Quan Lan, Jie-Xiong Mo, Wen-Biao Liu STOP** Recently, extended phase space thermodynamics of Lovelock AdS black holes has been of great interest. To provide

insight from (e) a different perspective and gain (e) a unified phase transition picture, non-extended phase space thermodynamics of (e) (n+1)-dimensional charged topological Lovelock AdS black holes is (=) investigated detailedly in (eb) the grand canonical ensemble. Specifically, the specific heat at (eb) constant electric potential is calculated and phase transition in the grand canonical ensemble is discussed. To probe the impact of the various parameters, they utilize (e) the control variate method and solve (e) the phase transition condition equation numerically for the case k=1, −1. There are two critical points for (e) the case n=6, k=1 while there is only one for other cases. For k=0, there exists no phase transition point. To figure out the nature of phase transition in the grand canonical ensemble, authors carry out an analytic check of the analog form of (e) Ehrenfest equations proposed by Banerjee et al. It is shown that Lovelock AdS black holes in (eb) the grand canonical ensemble undergo (e&eb) a second order phase transition. To examine the phase structure in (eb) the grand canonical ensemble, authors utilize the thermodynamic geometry method and calculate both the Weinhold metric and Ruppeiner metric. It is shown that for both analytic and graphical results that the divergence structure of (e) the Ruppeiner scalar curvature coincides with (=) that of the specific heat. Research provides (eb) one more example that Ruppeiner metric serves as (=) a wonderful tool to (e) probe the phase structures of (e) black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C75 (2015) 5, 211 DOI:10.1140/epjc/s10052-015-3434-5 Cite as: arXiv: 1503.01956 [gr-qc] (or arXiv: 1503.01956v1 [gr-qc] for this version **Non-extended phase space thermodynamics of Lovelock AdS black holes in grand canonical ensemble Jie-Xiong Mo, Wen-Biao Liu STOP** To provide an analytic verification of the nature of phase transition at the critical point of P−V criticality, the original expressions of Ehrenfest equations have been introduced directly. By treating the cosmological constant and (e&eb) its conjugate quantity as (=) thermodynamic pressure and volume respectively, authors carry out analytical check of (e) classical Ehrenfest equations. To show that approach is universal, authors investigate not only higher-dimensional charged AdS black holes (e&eb), but also rotating AdS black holes. Not only are the examples

of Einstein gravity shown, but also the example of modified gravity is presented for (e) Gauss-Bonnet AdS black holes. The specific heat at constant pressure CP, the volume expansion coefficient α and the isothermal compressibility coefficient κT are found to diverge exactly at the critical point. It has been verified that both Ehrenfest equations hold at (eb) the critical point of P−V criticality in (eb) the extended phase spaces of (e) AdS black holes. So the nature of the critical point of P−V criticality of AdS black holes has been demonstrated analytically to be (=) a second-order phase transition. These results are consistent with (e&eb) the nature of liquid-gas phase transition at the critical point. In this sense, extant research would deepen the understanding of (e) the relations of AdS black holes and (e&eb) liquid-gas systems. Moreover, successful approaches to introduce (eb) the original expressions of Erhenfest equations directly into (e&eb) black hole phase transition research demonstrate again (eb) that black hole thermodynamics is closely related (e&eb) to classical thermodynamics, which allows (eb) us to borrow techniques from (e) classical thermodynamics to investigate the thermodynamics of black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 89, 084057 (2014) DOI: 10.1103/PhysRevD.89.084057 Cite as:arXiv: 1404.3872 [gr-qc] (or arXiv:1404.3872v1 [gr-qc] for this version) **Ehrenfest scheme for P−V criticality of higher dimensional charged black holes, rotating black holes and Gauss-Bonnet AdS black holes Jie-Xiong Mo, Wen-Biao Liu STOP** Aiming at a unified phase transition picture of (e) the charged topological black hole in (eb) Ho\v{r}ava-Lifshitz gravity, authors investigate this issue not only in (eb) canonical ensemble with (e&eb) the fixed charge case but also in grand-canonical ensemble with (e&eb) the fixed potential case. Authors firstly perform the standard analysis of the specific heat,(e&eb) the free energy and the Gibbs potential, and (e&eb) then study its geometrothermodynamics. It is shown that the local phase transition points (eb) not only witness the divergence of (e) the specific heat, but also witness the minimum temperature and the maximum free energy or Gibbs potential. They also witness the divergence of (e) the corresponding thermodynamic scalar curvature. No matter which ensemble is chosen, the metric constructed can

successfully produce (eb) the behavior of the thermodynamic interaction and (e&eb) phase transition structure while other metrics failed to (e) predict the phase transition point of (e) the charged topological black hole in former literature. In grand-canonical ensemble, authors have discovered the phase transition which has not been reported before. It is similar to (e) the canonical ensemble in which the phase transition only takes place when k=−1. But it also has its unique characteristics that the location of the phase transition point depends on (eb) the value of potential, which is different from (e) the canonical ensemble where (e) the phase transition point is (=) independent of the parameters. After an analytical check of Ehrenfest scheme, authors find that the new phase transition is (=) a second order one. It is also found that the thermodynamics of (e) the black hole in Horava-Lifshitz gravity is quite different from that in (eb) Einstein gravity. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP1310(2013)056 DOI:10.1007/JHEP10(2013)056 Cite as: arXiv:1404.2497 [gr-qc] (or arXiv:1404.2497v1 [gr-qc] for this version) **A unified phase transition picture of the charged topological black hole in Horava-Lifshitz gravity Jie-Xiong Mo, Xiao-Xiong Zeng, Gu-Qiang Li, Xin Jiang, Wen-Biao Liu STOP** To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case $\beta\to\infty$ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness condition of (e) entropy. It is shown

that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure authors have discovered. It is shown that Legendre invariant thermodynamic scalar

curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws, leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions, geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged AdS black holes by considering the case of the RN-AdS black holes

surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase transition. For the case $\omega_q=-2/3$ authors derive analytic expressions of (e) critical physical quantities, while (e) for cases $\omega_q\neq-2/3$ appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this

version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=) chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy

Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed

quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions, while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma?**

All this consists of the act of consciousness alone", says Vasubandhu, "because unreal entities appear, just as a man with defective vision sees unreal hair or a moon, etc."He points to dreams as examples of purely subjective constructs which appear to be objective realities. The apparent reality dreams possess is not derived from any concrete, objective world, but merely from the idea of objectivity. While the Yogacara does not say that an idea has, for example, spatial attributes, it does have a form manifesting them. While he agrees with the Saiva idealist that

appearances have no independent existence apart from their appearing to consciousness, he maintains that for this reason they are unreal. The creativity of consciousness consists in its diversification in many modes having apparent externality; it is not a creation of objects. While the Kashmiri Saivite agrees that the world is pure conscious- ness alone, he maintains that it is such because it is a real creation of consciousness. The effect is essentially identical with the cause and shares in its reality. Matter and the entire universe are absolutely real, as 'congealed' (sty ana) or 'contracted' (samkucita) forms of consciousness. "This God of consciousness", writes Ksemaraja, "generates the universe and its form is a condensation of His own essence (rasa)" m By boiling sugarcane juice it condenses to form treacle, brown sugar and candy which retains its sweetness. Similarly, consciousness abides unchanged even though it assumes the concrete material form of the five gross elements. The same reality thus abides equally in gross and subtle forms. Consequently no object is totally insentient. Even stones bear a trace (vasana) of consciousness, although it is not clearly apparent because it is not associated with the vital breath (prana) and other components of a psycho-physical organism. Somananda goes so far as to affirm that physical objects, far from being insentient, can only exist insofar as they are aware of themselves as existing. The jar performs its function because it knows itself to be its agent. Indeed, all things are pervaded by consciousness and at one with it and hence share in its omniscience. Thus, Siva, Who perceives Himself in the form of physical objects, is the one ultimate reality. "The jar knows because it is of my nature", writes Somananda, "and I know it because I am of the jar's nature. I know because I am of Sadasiva's nature and He knows because He is of my nature; Yajnadatta [knows] because he is of Siva's nature and Siva [knows] because He is of Yajnadatta's nature". Integral Monism of Kashmiri Saivism Everything in this sense is directly perceived by absolute conscious- ness, and this direct perception (pratyaksa) unifies the knowable into a single, undivided whole. This is the central concept behind a doctrine originally expounded by Narasimha called 'the non-dualism of direct perception' (pratyaksadvaita). This states that consciousness is essentially perceptive and that its perception of all things operates throughout the universe. Insofar

as phenomena are clearly evident (sphufa) to us, everything is directly perceived by absolute consciousness, with which our individual consciousness is identical. This direct perception unfolds everywhere; the one true reality, it is alone and without companion or rival (nihsapatna). Even though it remains one, it can, by its very nature, perceive distinctions (bheda) between one entity and another, without this engendering any division within it. We distinguish between two entities in empirical terms on the basis of their mutual exclusion (anyonyabhava). The relative distinction {bheda) between them is essentially the perceived difference between their respective characteristics. Despite this difference they are united within the purview of a single cognition insofar as they are equally both manifest appearances. This cognition is the undivided essence (rasa) or 'own nature' (svabhava) of both. Encompassed by the 'fire of consciousness', there is no essential difference between them. Just as when an emerald and ruby reflect each other's light, the ruby is reddish-green and the emerald greenish-red, similarly everything is connected with everything else as part of the single variegated (vicitra) cognition of absolute consciousness. Mahesvarananda writes: The Supreme Lord's unique state of emotivity (asadharanabhava) is the outpouring of pure Being (mahasatta). It is manifest as the brilliance (sphuratta) of the universe which, if we ponder deeply, [is realized to be] the single flavour (ekarasa) of the essence of Beauty which is the vibration of the bliss of one's own nature. In this way all things are in reality one although divided from the one another sharing as they do the 'single flavour' (ekarasa) of the pure vibration of consciousness. Kashmiri Saiva Realism Kashmiri Saivism as a whole has been variously called a form of 'realistic idealism', 'monistic idealism', ‘idealistic monism' and 'concrete monism'. It is easy to understand why Kashmiri Saivism is The Doctrine of Vibration said to be 'idealistic' and 'monistic', but in what sense is it also 'realistic*? The answer to this question is of no small importance in trying to under- stand the central idea behind its metaphysics and the fundamental importance of the concept of Spanda, in this seemingly impossible marriage between monistic idealism and pluralistic realism. The Kashmiri Saiva approach understands the world to be a symbol of the absolute, that is, as the manner in which it presents itself to us. Again we can

contrast this view with that of the Advaita Vedanta. The Advaita Vedanta understands the world to be an expression of the absolute insofar as it exists by virtue of the absolute's Being. Being is understood to be the real unity which underlies empirically manifest separateness and as such is never empirically manifest. It is only transcendentally actual as 'being-in-itself. The Kashmiri Saiva position represents, in a sense, a reversal of this point of view. The nature of the absolute, and also that of Being, is onceived as an eternal becoming (satatodita), a dynamic flux or Spanda, 'the agency of the act of being'. It is dentified with the concrete actuality of the fact of appearing, not passive unmanifest Being. Appearance (abhasa) alone is real.ppearing (prakasamanatva) is equivalent to the fact of being (astitva). Ksemaraja writes in his commentary on the Stanzas on Vibration: Indeed, all things are manifest because they are nothing but manifestation. The point being that nothing is manifest apart from manifestation The absolutely unmanifest, from this point of view, can have as little existence as the space in a lattice window of a sky-palace. Nay, even less, because even that space can appear as an imagined mage manifest within consciousness. Everything is real according to the manner in which it appears. Even an illusion is in this sense real, insofar as it appears and is known in the manner in which it appears. The empirical and the real are identical categories of thought. As Abhinava says: Thus this is the supreme doctrine (upanisatf), namely that, when- ever and in whatever form [an entity] appears, that then is its particular nature. Perhaps at this stage a brief comparison with Heidegger's ideas might prove to be enlightening and not altogether out of place. According to Heidegger's phenomenology of Being, reality is intelligible in a two-fold Integral Monism of Kashmiri Saivism manner as 'phenomenon' and 'logos'. Heidegger defines what he means by 'phenomenon' as: "that-which-shows-itself. The manifest . . . phenomena are then the collection of that which lies open in broad daylight or can be brought to the light of day — what the Greeks at times implicitly identified as 'ta onta' (the things-which-are)".

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb) heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna**

**Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version) **Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb) strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel

period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar (=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI: 10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills (SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear

differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black**

**holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-

relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb)

the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this

version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstate thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11 August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of

entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.** Characterizing the behaviour of strongly coupled

quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at λ=λc, the correlation length diverges as ξ∼|λ−λc|−ν, while (e) the gap between the ground state and first excited state vanishes as (e) ΔE∼ξ−z∼|λ−λc|zν. Exactly at the critical point λ=λc, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as ω∼|k−k0|z. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-**

**equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004 STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element.

**Model has been applied for the following results:**

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| **Searches for Sterile Neutrinos with the IceCube Detector M. G. Aartsen et al. (IceCube Collaboration) Phys.**<br><br>**Rev. Lett 117, 071801 – Published 8 August 2016 Physics logo See Viewpoint: Hunting the Sterile Neutrino** |
| **(1) The IceCube neutrino telescope at the South Pole has measured the atmospheric muon neutrino spectrum as (=) a function of zenith angle and (e&eb) energy in the approximate 320 GeV to 20 TeV range, to search for (e) the oscillation signatures of (e) light sterile neutrinos.** |
| **(2) No evidence for anomalous νμ or ν¯μ disappearance is observed in (eb) either of two independently developed analyses, each using one year of (e) atmospheric neutrino data.** |
| **(3) New exclusion limits are placed on (e&eb) the parameter space of (e) the 3+1 model, in which muon antineutrinos experience (e) a strong Mikheyev-Smirnov-Wolfenstein-resonant oscillation.** |
| **(4) The exclusion limits extend to (e) sin22024≤0.02 at Δm2∼0.3 eV2 at the 90% confidence level.** |
| **(5) The allowed region from (e) global analysis of appearance experiments, including (e) LSND and MiniBooNE, is excluded at approximately the 99% confidence level for (e) the global best-fit value of \|Ue4\|2. Received 5 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.117.071801 © 2016 American Physical Society** |
| **Synopsis: Searching for Majorana Neutrinos August 16, 2016The KamLAND-Zen collaboration has run the most sensitive search to date for a radioactive decay that could reveal whether neutrinos are Majorana fermions. Synopsis figureKamLAND-Zen Collaboration** |
| **(6) Are neutrinos and (=) antineutrinos the same particles, as hypothesized by (e) physicist Ettore Majorana?** |
| **(7) If true, this "Majorana" nature of neutrinos could explain (eb) why neutrinos have (e) mass and why the Universe has (e) more matter than antimatter.** |
| **(8) Using an underground neutrino detector in Japan, the KamLAND-Zen collaboration has carried out the most sensitive search to date for a process called neutrinoless double-beta decay, which could confirm (eb) Majorana's hypothesis.** |
| **(9) The results set stringent constraints on the possible properties of (e) Majorana neutrinos.** |
| **(10) Double-beta decay is (=) a nuclear process in which two neutrons turn into (e&eb) two protons (or vice versa), emitting (eb) two antineutrinos.** |
| **(11) If neutrinos are Majorana fermions, (eb) the two antineutrinos could annihilate each other, and some double-beta decays would emit (eb) no antineutrinos.** |
| **(12) However, neutrinoless decay is believed to be (=) extremely rare. Searches of the process using (e) large amounts of the 35 naturally occurring (eb) isotopes that can undergo (e&eb) double-beta decay (such as xenon-136) have so far come up empty.** |

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| **(13) The KamLAND-Zen collaboration boosted their experimental sensitivity to (e) neutrinoless double-beta decays through (e&eb) a combination of factors.** |
| **(14) The researchers deployed clean detectors with (e&eb) very low background noise.** |
| **(15) And they used an unprecedented amount of xenon-136, which they purified to (e) remove any radioactive contaminants that could produce (eb) unwanted signals in the detectors.** |
| **(16) Combining data from several years, they improved (eb) the limit on the probability of neutrinoless double-beta decay by (e) sixfold compared to previous searches.** |
| **(17) This probability corresponds to (e&eb) a decay lifetime whose staggering value exceeds (e) 1026 years.** |
| **(18) This measurement allowed (eb) the researchers to set (eb) the most stringent upper limit on (e&eb) the possible mass of Majorana neutrinos (the particle must be lighter than 61–165 meV). This research is published in Physical Review Letters. –Matteo Rini Matteo Rini is the Deputy Editor of Physics.** |
| **How to Reach a Thousand-Second in-Plane Polarization Lifetime with 0.97−GeV/c Deuterons in a Storage Ring G. Guidoboni et al. (JEDI Collaboration) Phys. Rev. Lett. 117, 054801 – Published 28 July 2016 Physics logo See Synopsis: Deuterons Spin together for a Thousand Seconds** |
| **(19) Authors observe a deuteron beam polarization lifetime near 1000 s in (eb) the horizontal plane of (e) a magnetic storage ring (COSY).** |
| **(20) This long spin coherence time is maintained through (e&eb) a combination of beam bunching, electron cooling, (e&eb) sextupole field corrections, and (e&eb) the suppression of collective effects through (e&eb) beam current limits.** |
| **(21) This record lifetime is required for (e) a storage ring search for (e) an intrinsic electric dipole moment on (e&eb) the deuteron at a statistical sensitivity level approaching 10−29 e cm.** |
| **Synopsis: Deuterons Spin together for a Thousand Seconds July 28, 2016 Researchers set a new record for the in-plane spin-alignment lifetime of deuterons circulating in a magnetic storage ring. Synopsis figure E. Stephenson/Indiana University Center for Spacetime Symmetries Sime sentences and figures are deleted due to spatial restrictions.** |
| **(22) The electric dipole moment (EDM) of (e) a fundamental particle or a simple nucleus is believed to be (=) extremely small.** |
| **(23) One might therefore think that it isn't worth measuring this quantity. But it is. A nonzero EDM would signal (eb) a broken symmetry of (e) nature known as CP violation, which is required to (e) explain why the Universe has (e) more matter than antimatter.** |
| **(24) Edward Stephenson from Indiana University, Bloomington, and colleagues has now made a significant stride towards determining (eb) the EDM of the deuteron using (e) a deuteron beam circulating in (e&eb) a magnetic storage ring.** |

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| **(25) They have been able to preserve the alignment of the deuterons' spins for (e) a chunk of time that may be long enough to determine (eb) the deuteron's EDM.** |
| **(26) In experiments of this kind, the EDM can be determined by (e) measuring the rotation of (e) the initially aligned spins of (e) the beam particles as they travel around (e&eb) the ring.** |
| **(27) To obtain a measurement sufficiently accurate to explore (e&eb) CP violation predicted by (e) theories beyond the standard model, the spin alignment has to be (=) preserved for about 1000 s.** |
| **(28) But changes in particle speed—associated with (e&eb) position or direction variations among (e&eb) the particles in the beam—lead to (eb) a rapid loss of the alignment in the plane of the ring.** |
| **(29) Stephenson and co-workers have managed to retain (eb)the alignment of deuterons for the required 1000 s, which is 3 orders of magnitude longer (eb) than the previous record (obtained for an electron beam).** |
| **(30) Their success resulted from (e) capturing the beam into a tight bunch to reduce (e) position and direction spread and from (e) shaping the ring's magnetic field to(e) ensure all of the deuteron orbits around (e&eb) the ring have the same length, despite (e) direction variations. This research is published in Physical Review Letters. –Ana Lopes Ana Lopes is a Senior Editor of Physics.** |
| **Cotunneling Drag Effect in Coulomb-Coupled Quantum Dots A. J. Keller, J. S. Lim, David Sánchez, Rosa López, S. Amasha, J. A. Katine, Hadas Shtrikman, and D. Goldhaber-Gordon Phys. Rev. Lett. 117, 066602 – Published 4 August 2016 Physics logo See Synopsis: Coulomb Drag in a Double Dot** |
| **(31) In Coulomb drag, a current flowing in one conductor can induce (eb) a voltage across an adjacent conductor via (e&eb) the Coulomb interaction.** |
| **(32) The mechanisms yielding drag effects are not always understood, even though drag effects are (=) sufficiently general to be (=) seen in many low-dimensional systems.** |
| **(33) In this Letter, authors observe Coulomb drag in (eb) a Coulomb-coupled double quantum dot and, through (e&eb) both experimental and theoretical arguments, identify cotunneling as (=) essential to obtaining a correct qualitative understanding of the drag behavior.** |
| **Synopsis: Coulomb Drag in a Double Dot August 4, 2016 Electric current passing through a quantum dot can generate current in a nearby dot through a coordinated tunneling mechanism. Synopsis figureA. J. Keller et al., Phys. Rev. Lett. (2016)** |
| **(34) Electrons moving through (e&eb) a conductor can induce (eb) the movement of electrons of a nearby, but isolated, conductor.** |
| **(35) This so-called Coulomb drag is well understood in 2D layered materials, but the theory for nanowires and quantum dots does not (e) fully explain experiments.** |
| **(36) A new study of Coulomb drag in two nearby quantum dots shows (eb) that simultaneous tunneling of electrons—called cotunneling—can help explain the low-temperature behavior of (e) the system.** |

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| **(37) Electric current flows through (e&eb) a quantum dot by the tunneling of (e) electrons.** |
| **(38) When two dots are placed next to each other, the active, or lead, dot can drag current through (e&eb) the passive dot by (e) a process called sequential tunneling.** |
| **(39) In this case, an electron first tunnels onto (e&eb) the active dot, increasing (eb) the probability of tunneling onto (e&eb) the passive dot.** |
| **(40) Thus the passive dot follows (e&eb) the active dot's lead, but always one step behind.** |
| **(41) Sequential tunneling explains (eb) Coulomb drag observations at (eb) high temperature, but no studies have investigated cold temperatures where (e) quantum coherence effects may play a role.** |
| **(42) Andrew Keller from the California Institute of Technology, Pasadena, and his colleagues explored Coulomb drag at (eb) sub-Kelvin temperatures in a pair of quantum dots coupled by (e) capacitance.** |
| **(43) The team applied a current-driving voltage across (e&eb) one dot and detected a drag current in (eb) the other dot, even in regimes where (e) the driving voltage was very small.** |
| **(44) This conflicts with (e&eb) the sequential tunneling model, which predicts (eb) a threshold driving voltage below which (e) the drag current should be zero.** |
| **(45) The team showed that the inclusion of cotunneling into (eb) the drag theory explains the observations.** |
| **(46) Cotunneling could be important to efforts aimed at harvesting a unidirectional current from (e) a fluctuating voltage supply. This research is published in Physical Review Letters. –Michael Schirber Michael Schirber is a Corresponding Editor for Physics based in Lyon, France.** |
| **Maybe it wasn't the Higgs particle after all November 7, 2014 Maybe it wasn't the Higgs particle after all Associate Professor Mads Toudal Frandsen, University of Southern Denmark Credit: University of Southern Denmark** |
| **(47) Last year CERN announced the finding of a new elementary particle, the Higgs particle. But maybe it wasn't the Higgs particle, maybe it just looks like it. And maybe it is not alone. Many calculations indicate that the particle discovered last year in the CERN particle accelerator was (=) indeed the famous Higgs particle.** |
| **(48) Physicists agree that the CERN experiments did find a new particle that had never been seen before, but according to (e&eb) an international research team, there is (=) no conclusive evidence that the particle was indeed the Higgs particle.The research team has scrutinized the existing scientific data from CERN about the newfound particle and published their analysis in the journal Physical Review D. A member of this team is Mads Toudal Frandsen, associate professor at the Center for Cosmology and Particle Physics Phenomenology, Department of Physics, Chemistry and Pharmacy at the University of Southern Denmark. "The CERN data is generally taken as (=) evidence that the particle is the Higgs particle.** |

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| **(49) It is true that the Higgs particle can explain (eb) the data but there can be other explanations, we would also get this data from other particles", Mads Toudal Frandsen explains.** |
| **(50) The researchers' analysis does not (e) debunk the possibility that CERN has discovered the Higgs particle. That is still possible - but it is equally possible that it is a different kind of particle. "The current data is not precise enough to determine exactly what the particle is. It could be a number of other known particles", says Mads Toudal Frandsen. But if it wasn't the Higgs particle, which was found in CERN's particle accelerator, then what was it? "We believe that it may be a so-called techni-higgs particle. This particle is in some ways similar to the Higgs particle - hence half of the name", says Mads Toudal Frandsen.** |
| **(51) Although the techni-higgs particle and Higgs particle can easily be confused in experiments, they are two very different particles belonging to (e) two very different theories of how the universe was created.** |
| **(52) The Higgs particle is the missing piece in (eb) the theory called the Standard Model.** |
| **(53) This theory describes (eb) three of the four forces of nature. But it does not (e) explain what dark matter is - the substance that makes up (eb) most of the universe. A techni-higgs particle, if it exists, is (=) a completely different thing:** |
| **(54) "A techni-higgs particle is not (e) an elementary particle.** |
| **(55) Instead, it consists of so-called techni-quarks, which we believe are elementary. Techni-quarks may bind together in various ways to form (eb) for instance techni-higgs particles, while other combinations may form (eb) dark matter. Authors therefore expect to find several different particles at the LHC, all built by (e) techni-quarks", says Mads Toudal Frandsen.** |
| **(56) If techni-quarks exist, (eb) there must be a force to bind them together so that (e) they can form particles.** |
| **(57) None of the four known forces of nature (gravity, the electromagnetic force, the weak nuclear force and the strong nuclear force) are (=) any good at binding techni-quarks together.** |
| **(58) There must therefore be (=) a yet undiscovered force of nature. This force is called the the Technicolor force.** |
| **(59) What was found last year in CERN's accelerator could thus be (=) either the Higgs particle of the Standard Model or a light techni-higgs particle, composed of (e) two techni-quarks.** |
| **(60) Mads Toudal Frandsen believes that more data from CERN will probably be able to determine (eb) if it was a Higgs or a techni-higgs particle. If CERN gets an even more powerful accelerator, it will in principle be able to observe techni-quarks directly. More information: Technicolor Higgs boson in the light of LHC data. Phys. Rev. D 90, 035012th Alexander Belyaev, Matthew S. Brown, Roshan Foadi, and Mads T. Frandsen. journals.aps.org/prd/abstract/10.1103/PhysRevD.90.035012. On Arxiv: arxiv.org/abs/1309.2097 Journal reference: Physical Review D search and more info website: Provided by: University of Southern Denmark search and more info website** |

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| **The inflationary brane-antibrane universeClifford P. Burgess1,2, Mahbub Majumdar3, Detlef Nolte1, Fernando Quevedo3, Govindan Rajesh1 and Ren-Jie Zhang1 Published 13 August 2001 • Journal of High Energy Physics, Volume 2001, JHEP07(2001)** |
| **(61) Authors show how the motion through (e&eb) the extra dimensions of a gas of branes and antibranes can, under certain circumstances, produce (eb) an era of inflation as seen by (e) observers trapped on (e&eb) a 3-brane, with (e&eb) the inflaton being the inter-brane separation.** |
| **(62) Although most of discussion refers to (e) arbitrary ps, when we need to be specific they assume that they are (=) D-branes of type-II or type-I string theory.** |
| **(63) For realistic brane couplings, such as those arising in (eb) string theory, the inter-brane potentials are (=) too steep to inflate (eb) the universe for acceptably long times.** |
| **(64) However, for special regions of the parameter space of (e) brane-antibrane positions the brane motion is (=) slow enough for there to be (=) sufficient inflation.** |
| **(65) Inflation would be more generic in models where (e) the inter-brane interactions are much weaker.** |
| **(66) The spectrum of primordial density fluctuations predicted has (e) index n slightly less than 1, and acceptable amplitude, provided (e) that the extra dimensions have (e) linear size 1/r ~ 1012 GeV.** |
| **(67) Reheating occurs as in hybrid inflation, with (e&eb) the tachyonic instability of (e) the brane-antibrane system taking over for (e) small separations.** |
| **(68) The tachyon field can induce (eb) a cascade mechanism within which higher-dimension branes annihilate (e) into lower-dimension ones.** |
| **(69) They argue that such a cascade naturally stops with (e&eb) the production of 3-branes in (eb) 10-dimensional string theory.** |
| **Effect of the neutron lifetime on processes in the early universe A. V. Chechkin, A. V. Ivanchik, A. P. Serebrov, S. V. Bobashev First Online: 27 July 2016 DOI: 10.1134/S1063784216070069 Chechkin, A.V., Ivanchik, A.V., Serebrov, A.P. et al. Tech. Phys. (2016) 61: 1101. Doi: 10.1134/S1063784216070069** |
| **(70) The influence of the neutron lifetime on (e&eb) the abundance of light elements produced during (e&eb) the primordial nucleosynthesis early in (eb) the birth of the Universe is considered.** |
| **(71) Among light elements, namely, D, 3He, 4He, and 7Li, 4He proves to be (=) most sensitive to neutron lifetime τn.** |
| **(72) Astronomic data on the light element abundance also provide (eb) the best accuracy for 4He.** |
| **(73) The solution of a number of problems discussed in this paper requires (e) improving the accuracy of observations for the 4He abundance and refining the value of τn.** |

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| **Application of a Theory and Simulation based Convective Boundary Mixing model for AGB Star Evolution and Nucleosynthesis U. Battino, M. Pignatari, C. Ritter, F. Herwig, P. Denisenkov, J.W. Den Hartogh, R. Trappitsch, R. Hirschi, B. Freytag, F. Thielemann, B. Paxton (Submitted on 19 May 2016)** |
| **(74) The s-process nucleosynthesis in Asymptotic Giant Branch (AGB) stars depends on (e&eb) the modeling of (e) convective boundaries.** |
| **(75) Authors' present models and s-process simulations that adopt (e&eb) a treatment of convective boundaries based on (e) the results of hydrodynamic simulations and on (e&eb) the theory of mixing due to (e) gravity waves in (eb) the vicinity of convective boundaries.** |
| **(76) Hydrodynamics simulations suggest (eb) the presence of convective boundary mixing (CBM) at (eb) the bottom of the thermal pulse-driven convective zone.** |
| **(77) Similarly, convection-induced mixing processes are proposed for (e) the mixing below the convective envelope during (e&eb) third dredge-up where (e) the 13C pocket for (e) the s process in AGB stars forms.** |
| **(78) In this work apply apply a CBM model motivated by (e) simulations and theory to (e&eb) models with initial mass M = 2 and M = 3M?, and with (e&eb) initial metal content Z = 0:01 and Z = 0:02.** |
| **(79) As reported previously, the He-intershell abundance of 12C and 16O are increased by (e) CBM at the bottom of (e) pulse-driven convection zone.** |
| **(80) This mixing is affecting (e&eb) the 22Ne (alpha, n) 25Mg activation and (e&eb) the s-process effciency in the 13C-pocket.** |
| **(81) In our model CBM at the bottom of the convective envelope during (e&eb) the third dredgeup represents (eb) gravity wave mixing.** |
| **(82) Authors take further into account that hydrodynamic simulations indicate (eb) a declining mixing efficiency already about (e&eb) a pressure scale height from (e) the convective boundaries, compared to (e&eb) mixing-length theory.** |
| **(83) The final s-process abundances are characterized by (e) 0:36 < [s=Fe] < 0:78 and the heavy-to-light s-process ratio is (=) 0:23 < [hs=ls] < 0:45.** |
| **(84) Finally, authors compare results with (e&eb) stellar observations, pre-solar grain measurements and previous work. Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI:10.3847/0004-637X/827/1/30 Cite as: arXiv:1605.06159 [astro-ph.SR] (or arXiv:1605.06159v1 [astro-ph.SR] for this version)** |
| **MESA and NuGrid simulations of classical novae: CO and ONe nova nucleosynthesis Pavel A. Denissenkov (UVic), James W. Truran (Chicago), Marco Pignatari (Basel), Reto Trappitsch (Chicago), Christian Ritter (Frankfurt), Falk Herwig (UVic), Umberto Battino (Basel), Kiana Setoodehnia (McMaster), B. Paxton (UC Santa Barbara)** |

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| **(85) Classical novae are the result of (e) thermonuclear flashes of hydrogen accreted by (e) CO or ONe white dwarfs, leading (eb) eventually to the dynamic ejection of (e) the surface layers.** |
| **(86) These are observationally known to be (=) enriched in heavy elements, such as (=) C, O and Ne that must originate in (eb) layers below the H-flash convection zone.** |
| **(87) Building on their previous work, they now present stellar evolution simulations of (e) ONe novae and provide (eb) a comprehensive comparison of models with (e&eb) published ones.** |
| **(88) Some of authors models include exponential convective boundary mixing to account for (e) the observed enrichment of (e) the nova ejecta even when accreted material has (e) a solar abundance distribution.** |
| **(89) Models produce (eb) maximum temperature evolution profiles and nucleosynthesis yields (eb) in good agreement with (e&eb) models that generate (eb) enriched ejecta by assuming (eb) that the accreted material was pre-mixed.** |
| **(90) Authors confirm for (e) ONe novae the result we reported previously, i.e.\ we found that 3He could be produced {\it in situ} in (eb) solar-composition envelopes accreted with (e&eb) slow rates ($M˙$<10−10$M_\odot$/yr) by (e) cold (TWD<107 K) CO WDs, and that convection was triggered by (eb) 3He burning before (e) the nova outburst in that case.** |
| **(91) In addition, authors now find that the interplay between the 3He production and (e&eb) destruction in the solar-composition envelope accreted with (e&eb) an intermediate rate, e.g.\ $M˙$=10−10$M_\odot$/yr, by the 1.15$M_\odot$ ONe WD with (e&eb) a relatively high initial central temperature, e.g.\ TWD=15×106 K, leads to (eb) the formation of a thick radiative buffer zone that separates (e&eb) the bottom of the convective envelope from (e) the WD surface. (Abridged) Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI: 10.1093/mnras/stu1000 Cite as: arXiv:1303.6265 [astro-ph.SR] (or arXiv:1303.6265v2 [astro-ph.SR] for this version)** |
| **Hybrid C-O-Ne White Dwarfs as Progenitors of Diverse SNe Ia Pavel Denissenkov (UVic), James Truran (U. Chicago), Falk Herwig (UVic), Sam Jones (UVic), Bill Paxton (UC Santa Barbara), Ken'ichi Nomoto (U. Tokyo), Toshio Suzuki (Nihon U.), Hiroshi Toki (Osaka U.)(Submitted on 6 Nov 2014)** |
| **(92) When carbon is ignited off-center in a CO core of a super-AGB star (eb) its burning in a convective shell tends to propagate to (e&eb) the center.** |
| **(93) Whether the C flame will actually be able to reach (eb) the center depends on (e&eb) the efficiency of (e) extra mixing beneath (e&eb) the C convective shell.** |
| **(94) Whereas thermohaline mixing is too inefficient to interfere with the C-flame propagation, (eb) convective boundary mixing can prevent (e) the C burning from reaching (eb) the center.** |
| **(95) As a result, a C-O-Ne white dwarf (WD) is (=) formed, after the star has lost (e) its envelope.** |
| **(96) Such a "hybrid" WD has (e) a small CO core surrounded by (e) a thick ONe zone.** |

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| **(97) In the proposed 1D stellar evolution computations the hybrid WD is allowed to (eb) accrete C-rich material, as if it were in a close binary system and (e&eb) accreted H-rich material from (e) its companion with (e&eb) a sufficiently high rate at which the accreted H would be processed into (e&eb) He under (e&eb) stationary conditions, assuming (eb) that He could then be transformed into (e&eb) C.** |
| **(98) When the mass of the accreting WD approaches (eb) the Chandrasekhar limit, they find a series of convective Urca shell flashes associated with (e&eb) high abundances of 23Na and 25Mg.** |
| **(99) They are followed by off-center C ignition leading to (eb) convection that occupies (e) almost the entire star.** |
| **(100) To model the Urca processes, authors use the most recent well-resolved data for (e) their reaction and (e&eb) neutrino-energy loss rates.** |
| **(101) Because of the emphasized uncertainty of the convective Urca process in hybrid WD models of (e) SN Ia progenitors, authors consider a number of their potentially possible alternative instances for (e) different mixing assumptions, all of which reach (eb) a phase of explosive C ignition, either off or in (eb) the center.** |
| **(102) Our hybrid SN Ia progenitor models have (e) much lower C to O abundance ratios at (eb) the moment of the explosive C ignition than (e) their pure CO counterparts, which may explain (eb) the observed diversity of the SNe Ia. subjects: Solar and Stellar Astrophysics (astro-ph.SR)cite as: arXiv:1411.1471 [astro-ph.SR](or arXiv:1411.1471v1 [astro-ph.SR] for this version)** |
| **Rejecting proposed dense-matter equations of state with quiescent low-mass X-ray binaries Sebastien Guillot, Robert E. Rutledge (Submitted on 15 Sep 2014 (v1), last revised 20 Oct 2014 (this version, v2))** |
| **(103) Neutrons stars are unique laboratories to (e) discriminate between the various proposed equations of (e) state of matter at and above nuclear density.** |
| **(104) One sub-class of neutron stars - those inside (eb) quiescent low-mass X-ray binaries (qLMXBs) - produce (eb) a thermal surface emission from (e) which the neutron star radius (R_NS) can be measured, using (e) the widely accepted observational scenario for qLMXBs, assuming (eb) unmagnetized H atmospheres.** |
| **(105) In a combined spectral analysis, this work first reproduces (eb) a previously published measurement of the \rns, assumed to be (=) the same for all neutron stars, using (e) a slightly expanded data set.** |
| **(106) The radius measured is R_NS = 9.4 +/-1.2 km. On the basis of spectral analysis alone, this measured value is not affected by imposing (e&eb) an assumption of causality in the core.** |
| **(107) However, the assumptions underlying this R_NS measurement would be falsified by (e) the observation of any neutron star with (e&eb) a mass >2.6 Msun, since radii <11 km would be rejected if (e) causality is assumed, which would exclude (eb) most of the R_NS parameter space obtained in this analysis.** |

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| **(108) Finally, this work directly tests (e&eb) a selection of dense matter equations of states: WFF1, AP4, MPA1, PAL1, MS0, and three versions of equations of state produced through (e&eb) chiral effective theory.** |
| **(109) Two of those, MSO and PAL1, are rejected at (eb) the 99% confidence level, accounting for (e) all quantifiable uncertainties, while (e) the other cannot be excluded at >99% certainty. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE) DOI: 10.1088/2041-8205/796/1/L3 Cite as: arXiv:1409.4306 [astro-ph.HE] (or arXiv:1409.4306v2 [astro-ph.HE] for this version** |

## SECTION FIVE

## Photon-Recoil Heating Rate For Various Experimental Parameters And Quantum Predictions

**INTRODUCTION—VARIABLES USED**

**Viewpoint: Measuring Quantum Kicks from a Beam of Light Amir Safavi-Naeini, 348 Via Pueblo Mall, Stanford University, Stanford, CA 94305-4088, USA June 13, 2016• Physics 9, 64 Force sensors levitated by light have reached the quantum regime, in which their sensitivity is limited by the momentum kicks of individual photons. Figure caption expand figure APS/Alan Stonebraker Figures and some sentences are deleted due to spatial restrictions Please excuse me**

(1) They then measured the rate at which this warm-up occurred, over many runs and at ultralow pressures, and determined that its value was on (e&eb) the order of a few tens of kilohertz—an impressively slow rate for (e) a room-temperature experiment.
(2) In addition, the measured values of the photon-recoil heating rate for various experimental parameters agree with (=) quantum predictions.
(3) Since the first proposals for using (e) trapped nanoparticles in (eb) quantum optical and optomechanical experiments [9, 10], major efforts have been directed at using (e) these systems to observe (e) quantum phenomena.
(4) Such sensors might allow us to someday test quantum mechanics itself in the limit of (e) "massive" systems, which, like the nanosphere, consist of (e) hundreds of millions of atoms [11].
(5) By demonstrating a measurement of photon recoil heating, Novotny's group has unambiguously observed (e) the first signatures of quantum phenomena on (e&eb) optically levitated nanoparticles

**NOTATION**

**Module One**

To measure the photon-recoil heating rate, the ETH group turned off their mechanism for cooling the particle by holding the laser intensity constant and allowing the particle to heat up through its interaction with (e&eb) the (room-temperature) environment

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| $G_{13}$ : Category one of **measure the photon-recoil heating rate, the ETH group turned off their mechanism for cooling the particle by holding the laser intensity constant and allowing the particle to heat up through its interaction;** (room-temperature) environment<br><br>$G_{14}$ : Category two of SAS<br><br>$G_{15}$ : Category three of SAS | |
| $T_{13}$ : Category one of (room-temperature) environment **; measure the photon-recoil heating rate, the ETH group turned off their mechanism for cooling the particle by holding the laser intensity constant and allowing the particle to heat up through its interaction**<br><br>$T_{14}$ : Category two of SAS<br><br>$T_{15}$ : Category three of SAS | |
| **Module Two**<br><br>They then measured the rate at which this warm-up occurred, over many runs and at ultralow pressures, and determined that its value was of the order of a few tens of kilohertz—an impressively slow rate for a room-temperature experiment | |
| $G_{16}$ : Category one of measured the rate at which this warm-up occurred, over many runs and at ultralow pressures, and determined that its value<br><br>$G_{17}$ : Category two of SAS<br><br>$G_{18}$ : Category three of SAS | |
| $T_{16}$ : Category one of order of a few tens of kilohertz—an impressively slow rate for (e) a room-temperature experiment<br><br>$T_{17}$ : Category two of SAS<br><br>$T_{18}$ : Category three of SAS | |
| **Module three**<br><br>In addition, the measured values of the photon-recoil heating rate for various experimental parameters agree with (=) quantum predictions | |
| $G_{20}$ : Category one of measured values of the photon-recoil heating rate for various experimental parameters<br><br>$G_{21}$ : Category two of SAS<br><br>$G_{22}$ : Category three of SAS | |
| $T_{20}$ : Category one of quantum predictions<br><br>$T_{21}$ : Category two of SAS | |

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| $T_{22}$ : Category three of SAS | |
| **Module four**<br><br>Since the first proposals for using trapped nanoparticles in (eb) quantum optical and optomechanical experiments [9, 10], major efforts have been directed at using (e) these systems to observe (e) quantum phenomena | |
| $G_{24}$ : Category one of trapped nanoparticles<br><br>$G_{25}$ : Category two of SAS<br><br>$G_{26}$ : Category three of SAS | |
| $T_{24}$ : Category one of quantum optical and optomechanical experiments [9, 10], major efforts have been directed at using (e) these systems to observe (e) quantum phenomena<br><br>$T_{25}$ : Category two of SAS<br><br>$T_{26}$ : Category three of SAS | |
| **Module five**<br><br>Since the first proposals for using trapped nanoparticles in quantum optical and optomechanical experiments [9, 10], major efforts have been directed at using (e) these systems to observe (e) quantum phenomena | |
| $G_{28}$ : Category one of these systems to observe (e) quantum phenomena<br><br>$G_{29}$ : Category two of SAS<br><br>$G_{30}$ : Category three of SAS | |
| $T_{28}$ : Category one of proposals for using trapped nanoparticles in quantum optical and optomechanical experiments [9, 10], major efforts have been directed<br><br>$T_{29}$ : Category two of SAS<br><br>$T_{30}$ : Category three of SAS | |
| **Module six**<br><br>Since the first proposals for using trapped nanoparticles in quantum optical and optomechanical experiments [9, 10], major efforts have been directed at using (e) these systems to observe (e) quantum phenomena | |
| $G_{32}$ : Category one of these systems to observe (e) quantum phenomena<br><br>$G_{33}$ : Category two of SAS<br><br>$G_{34}$ : Category three of SAS | |

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| $T_{32}$ : Category one of proposals for using trapped nanoparticles in quantum optical and optomechanical experiments [9, 10], major efforts have been directed<br><br>$T_{33}$ : Category two of SAS<br><br>$T_{34}$ : Category three of SAS | |
| **Module seven**<br><br>Since the first proposals for using trapped nanoparticles in quantum optical and optomechanical experiments [9, 10], major efforts have been directed at using these systems to observe (e) quantum phenomena | |
| $G_{36}$ : Category one of quantum phenomena<br><br>$G_{37}$ : Category two of SAS<br><br>$G_{38}$ : Category three of SAS | |
| $T_{36}$ : Category one of first proposals for using trapped nanoparticles in quantum optical and optomechanical experiments [9, 10], major efforts have been directed at using these systems<br><br>$T_{37}$ : Category two of SAS<br><br>$T_{38}$ : Category three of SAS | |
| **Module eight**<br><br>Such sensors might allow us to someday test quantum mechanics itself in the limit of (e) "massive" systems, which, like the nanosphere, consist of (e) hundreds of millions of atoms [11] | |
| $G_{40}$ : Category one of"massive" systems, which, like the nanosphere, consist of (e) hundreds of millions of atoms [11]<br><br>$G_{41}$ : Category two of SAS<br><br>$G_{42}$ : Category three of SAS | |
| $T_{40}$ : Category one of Such sensors might allow us to someday test quantum mechanics itself in the limit<br><br>$T_{41}$ : Category two of SAS<br><br>$T_{42}$ : Category three of SAS | |
| **Module Nine**<br><br>Such sensors might allow us to someday test quantum mechanics itself in the limit of "massive" systems, which, like the nanosphere, consist of (e) hundreds of millions of atoms [11] | |
| $G_{44}$ : Category one of hundreds of millions of atoms [11] | |

$G_{45}$ : Category two of SAS

$G_{46}$ : Category three of SAS

$T_{44}$ : Category one of sensors might allow us to someday test quantum mechanics itself in the limit of "massive" systems, which, like the nanosphere

$T_{45}$ : Category two of SAS

$T_{46}$ : Category three of SAS

**The Coefficients:**

$(a_{13})^{(1)}, (a_{14})^{(1)}, (a_{15})^{(1)}, (b_{13})^{(1)}, (b_{14})^{(1)}, (b_{15})^{(1)}\ (a_{16})^{(2)}, (a_{17})^{(2)}, (a_{18})^{(2)}\ (b_{16})^{(2)}, (b_{17})^{(2)}, (b_{18})^{(2)}$:
$(a_{20})^{(3)}, (a_{21})^{(3)}, (a_{22})^{(3)}$ , $(b_{20})^{(3)}, (b_{21})^{(3)}, (b_{22})^{(3)}$
$(a_{24})^{(4)}, (a_{25})^{(4)}, (a_{26})^{(4)}, (b_{24})^{(4)}, (b_{25})^{(4)}, (b_{26})^{(4)}, (b_{28})^{(5)}, (b_{29})^{(5)}, (b_{30})^{(5)}, (a_{28})^{(5)}, (a_{29})^{(5)}, (a_{30})^{(5)}$,
$(a_{32})^{(6)}, (a_{33})^{(6)}, (a_{34})^{(6)}, (b_{32})^{(6)}, (b_{33})^{(6)}, (b_{34})^{(6)}$
$(a_{36})^{(7)}, (a_{37})^{(7)}, (a_{38})^{(7)}, (b_{36})^{(7)}, (b_{37})^{(7)}, (b_{38})^{(7)}$
$(a_{40})^{(8)}, (a_{41})^{(8)}, (a_{42})^{(8)}, (b_{40})^{(8)}, (b_{41})^{(8)}, (b_{42})^{(8)}$
$(a_{44})^{(9)}, (a_{45})^{(9)}, (a_{46})^{(9)}, (b_{44})^{(9)}, (b_{45})^{(9)}, (b_{46})^{(9)}$

are Accentuation coefficients
$(a_{13}^{'})^{(1)}, (a_{14}^{'})^{(1)}, (a_{15}^{'})^{(1)}, (b_{13}^{'})^{(1)}, (b_{14}^{'})^{(1)}, (b_{15}^{'})^{(1)}, (a_{16}^{'})^{(2)}, (a_{17}^{'})^{(2)}, (a_{18}^{'})^{(2)}, \ (b_{16}^{'})^{(2)}, (b_{17}^{'})^{(2)}, (b_{18}^{'})^{(2)}$
$, (a_{20}^{'})^{(3)}, (a_{21}^{'})^{(3)}, (a_{22}^{'})^{(3)}, (b_{20}^{'})^{(3)}, (b_{21}^{'})^{(3)}, (b_{22}^{'})^{(3)}$
$(a_{24}^{'})^{(4)}, (a_{25}^{'})^{(4)}, (a_{26}^{'})^{(4)}, (b_{24}^{'})^{(4)}, (b_{25}^{'})^{(4)}, (b_{26}^{'})^{(4)}, (b_{28}^{'})^{(5)}, (b_{29}^{'})^{(5)}, (b_{30}^{'})^{(5)}\ (a_{28}^{'})^{(5)}, (a_{29}^{'})^{(5)}, (a_{30}^{'})^{(5)}$ ,
$(a_{32}^{'})^{(6)}, (a_{33}^{'})^{(6)}, (a_{34}^{'})^{(6)}, (b_{32}^{'})^{(6)}, (b_{33}^{'})^{(6)}, (b_{34}^{'})^{(6)}$
$(a_{36}^{'})^{(7)}, (a_{37}^{'})^{(7)}, (a_{38}^{'})^{(7)}, (b_{36}^{'})^{(7)}, (b_{37}^{'})^{(7)}, (b_{38}^{'})^{(7)}$,
$(a_{40}^{'})^{(8)}, (a_{41}^{'})^{(8)}, (a_{42}^{'})^{(8)}, (b_{40}^{'})^{(8)}, (b_{41}^{'})^{(8)}, (b_{42}^{'})^{(8)}$,
$(a_{44}^{'})^{(9)}, (a_{45}^{'})^{(9)}, (a_{46}^{'})^{(9)}, (b_{44}^{'})^{(9)}, (b_{45}^{'})^{(9)}, (b_{46}^{'})^{(9)}$,
are Dissipation coefficients

**Module Numbered One**

**The differential system of this model is now (Module Numbered one)**

$$\frac{dG_{13}}{dt} = (a_{13})^{(1)}G_{14} - \left[(a_{13}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14}, t)\right]G_{13} \quad (1)$$

$$\frac{dG_{14}}{dt} = (a_{14})^{(1)}G_{13} - \left[(a_{14}^{'})^{(1)} + (a_{14}^{''})^{(1)}(T_{14}, t)\right]G_{14} \quad (2)$$

$$\frac{dG_{15}}{dt} = (a_{15})^{(1)}G_{14} - \left[(a_{15}^{'})^{(1)} + (a_{15}^{''})^{(1)}(T_{14}, t)\right]G_{15} \quad (3)$$

$$\frac{dT_{13}}{dt} = (b_{13})^{(1)}T_{14} - \left[(b_{13}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G, t)\right]T_{13} \quad (4)$$

$$\frac{dT_{14}}{dt} = (b_{14})^{(1)}T_{13} - \left[(b_{14}^{'})^{(1)} - (b_{14}^{''})^{(1)}(G, t)\right]T_{14} \quad (5)$$

$$\frac{dT_{15}}{dt} = (b_{15})^{(1)}T_{14} - \left[(b_{15}^{'})^{(1)} - (b_{15}^{''})^{(1)}(G, t)\right]T_{15} \quad (6)$$

$+(a_{13}^{''})^{(1)}(T_{14}, t) =$ First augmentation factor

$-(b_{13}^{''})^{(1)}(G, t) =$ First detritions factor

**Module Numbered Two**

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| **The differential system of this model is now ( Module numbered two)** | |
| $\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \left[(a_{16}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17},t)\right]G_{16}$ | 7 |
| $\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \left[(a_{17}^{'})^{(2)} + (a_{17}^{''})^{(2)}(T_{17},t)\right]G_{17}$ | 8 |
| $\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \left[(a_{18}^{'})^{(2)} + (a_{18}^{''})^{(2)}(T_{17},t)\right]G_{18}$ | 9 |
| $\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \left[(b_{16}^{'})^{(2)} - (b_{16}^{''})^{(2)}((G_{19}),t)\right]T_{16}$ | 10 |
| $\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \left[(b_{17}^{'})^{(2)} - (b_{17}^{''})^{(2)}((G_{19}),t)\right]T_{17}$ | 11 |
| $\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \left[(b_{18}^{'})^{(2)} - (b_{18}^{''})^{(2)}((G_{19}),t)\right]T_{18}$ | 12 |
| $+(a_{16}^{''})^{(2)}(T_{17},t) =$ First augmentation factor | |
| $-(b_{16}^{''})^{(2)}((G_{19}),t) =$ First detritions factor | |
| **Module Numbered Three**<br><br>**The differential system of this model is now (Module numbered three)** | |
| $\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \left[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21},t)\right]G_{20}$ | 13 |
| $\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \left[(a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)}(T_{21},t)\right]G_{21}$ | 14 |
| $\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \left[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21},t)\right]G_{22}$ | 15 |
| $\frac{dT_{20}}{dt} = (b_{20})^{(3)}T_{21} - \left[(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23},t)\right]T_{20}$ | 16 |
| $\frac{dT_{21}}{dt} = (b_{21})^{(3)}T_{20} - \left[(b_{21}^{'})^{(3)} - (b_{21}^{''})^{(3)}(G_{23},t)\right]T_{21}$ | 17 |
| $\frac{dT_{22}}{dt} = (b_{22})^{(3)}T_{21} - \left[(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23},t)\right]T_{22}$ | 18 |
| $+(a_{20}^{''})^{(3)}(T_{21},t) =$ First augmentation factor | |
| $-(b_{20}^{''})^{(3)}(G_{23},t) =$ First detritions factor | |
| **Module Numbered Four**<br><br>**The differential system of this model is now (Module numbered Four)** | |
| $\frac{dG_{24}}{dt} = (a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25},t)\right]G_{24}$ | 19 |
| $\frac{dG_{25}}{dt} = (a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25},t)\right]G_{25}$ | 20 |
| $\frac{dG_{26}}{dt} = (a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25},t)\right]G_{26}$ | 21 |
| $\frac{dT_{24}}{dt} = (b_{24})^{(4)}T_{25} - \left[(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}((G_{27}),t)\right]T_{24}$ | 22 |
| $\frac{dT_{25}}{dt} = (b_{25})^{(4)}T_{24} - \left[(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}((G_{27}),t)\right]T_{25}$ | 23 |
| $\frac{dT_{26}}{dt} = (b_{26})^{(4)}T_{25} - \left[(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}((G_{27}),t)\right]T_{26}$ | 24 |
| $+(a_{24}^{''})^{(4)}(T_{25},t) =$ **First augmentation factor** | |
| $-(b_{24}^{''})^{(4)}((G_{27}),t) =$ **First detritions factor** | |
| **Module Numbered Five**:<br><br>**The differential system of this model is now (Module number five)** | |
| $\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29},t)\right]G_{28}$ | 25 |
| $\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29},t)\right]G_{29}$ | 26 |

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| $\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \left[(a'_{30})^{(5)} + (a''_{30})^{(5)}(T_{29}, t)\right]G_{30}$ | 27 |
| $\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \left[(b'_{28})^{(5)} - (b''_{28})^{(5)}((G_{31}), t)\right]T_{28}$ | 28 |
| $\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \left[(b'_{29})^{(5)} - (b''_{29})^{(5)}((G_{31}), t)\right]T_{29}$ | 29 |
| $\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \left[(b'_{30})^{(5)} - (b''_{30})^{(5)}((G_{31}), t)\right]T_{30}$ | 30 |
| $+(a''_{28})^{(5)}(T_{29}, t) =$ **First augmentation factor** | |
| $-(b''_{28})^{(5)}((G_{31}), t) =$ **First detritions factor** | |
| **Module Numbered Six**<br><br>**The differential system of this model is now (Module numbered Six)** | |
| $\frac{dG_{32}}{dt} = (a_{32})^{(6)}G_{33} - \left[(a'_{32})^{(6)} + (a''_{32})^{(6)}(T_{33}, t)\right]G_{32}$ | 31 |
| $\frac{dG_{33}}{dt} = (a_{33})^{(6)}G_{32} - \left[(a'_{33})^{(6)} + (a''_{33})^{(6)}(T_{33}, t)\right]G_{33}$ | 32 |
| $\frac{dG_{34}}{dt} = (a_{34})^{(6)}G_{33} - \left[(a'_{34})^{(6)} + (a''_{34})^{(6)}(T_{33}, t)\right]G_{34}$ | 33 |
| $\frac{dT_{32}}{dt} = (b_{32})^{(6)}T_{33} - \left[(b'_{32})^{(6)} - (b''_{32})^{(6)}((G_{35}), t)\right]T_{32}$ | 34 |
| $\frac{dT_{33}}{dt} = (b_{33})^{(6)}T_{32} - \left[(b'_{33})^{(6)} - (b''_{33})^{(6)}((G_{35}), t)\right]T_{33}$ | 35 |
| $\frac{dT_{34}}{dt} = (b_{34})^{(6)}T_{33} - \left[(b'_{34})^{(6)} - (b''_{34})^{(6)}((G_{35}), t)\right]T_{34}$ | 36 |
| $+(a''_{32})^{(6)}(T_{33}, t) =$ **First augmentation factor** | |
| **Module Numbered Seven:**<br><br>**The differential system of this model is now (Seventh Module)** | |
| $\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \left[(a'_{36})^{(7)} + (a''_{36})^{(7)}(T_{37}, t)\right]G_{36}$ | 37 |
| $\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \left[(a'_{37})^{(7)} + (a''_{37})^{(7)}(T_{37}, t)\right]G_{37}$ | *38* |
| $\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \left[(a'_{38})^{(7)} + (a''_{38})^{(7)}(T_{37}, t)\right]G_{38}$ | *39* |
| $\frac{dT_{36}}{dt} = (b_{36})^{(7)}T_{37} - \left[(b'_{36})^{(7)} - (b''_{36})^{(7)}((G_{39}), t)\right]T_{36}$ | *40* |
| $\frac{dT_{37}}{dt} = (b_{37})^{(7)}T_{36} - \left[(b'_{37})^{(7)} - (b''_{37})^{(7)}((G_{39}), t)\right]T_{37}$ | *41* |
| $\frac{dT_{38}}{dt} = (b_{38})^{(7)}T_{37} - \left[(b'_{38})^{(7)} - (b''_{38})^{(7)}((G_{39}), t)\right]T_{38}$ | *42* |
| $+(a''_{36})^{(7)}(T_{37}, t) =$ **First augmentation factor** | |
| **Module Numbered Eight**<br><br>The differential system of this model is now | |
| $\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \left[(a'_{40})^{(8)} + (a''_{40})^{(8)}(T_{41}, t)\right]G_{40}$ | *43* |
| $\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \left[(a'_{41})^{(8)} + (a''_{41})^{(8)}(T_{41}, t)\right]G_{41}$ | *44* |
| $\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \left[(a'_{42})^{(8)} + (a''_{42})^{(8)}(T_{41}, t)\right]G_{42}$ | *45* |
| $\frac{dT_{40}}{dt} = (b_{40})^{(8)}T_{41} - \left[(b'_{40})^{(8)} - (b''_{40})^{(8)}((G_{43}), t)\right]T_{40}$ | *46* |
| $\frac{dT_{41}}{dt} = (b_{41})^{(8)}T_{40} - \left[(b'_{41})^{(8)} - (b''_{41})^{(8)}((G_{43}), t)\right]T_{41}$ | *47* |
| $\frac{dT_{42}}{dt} = (b_{42})^{(8)}T_{41} - \left[(b'_{42})^{(8)} - (b''_{42})^{(8)}((G_{43}), t)\right]T_{42}$ | *48* |
| **Module Numbered Nine** | |

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| The differential system of this model is now | |
| $\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \left[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45},t)\right]G_{44}$ | 49 |
| $\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \left[(a_{45}^{'})^{(9)} + (a_{45}^{''})^{(9)}(T_{45},t)\right]G_{45}$ | 50 |
| $\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \left[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45},t)\right]G_{46}$ | 51 |
| $\frac{dT_{44}}{dt} = (b_{44})^{(9)}T_{45} - \left[(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}((G_{47}),t)\right]T_{44}$ | 52 |
| $\frac{dT_{45}}{dt} = (b_{45})^{(9)}T_{44} - \left[(b_{45}^{'})^{(9)} - (b_{45}^{''})^{(9)}((G_{47}),t)\right]T_{45}$ | 53 |
| $\frac{dT_{46}}{dt} = (b_{46})^{(9)}T_{45} - \left[(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}((G_{47}),t)\right]T_{46}$ | 54 |
| $+(a_{44}^{''})^{(9)}(T_{45},t) =$ **First augmentation factor** | |
| $-(b_{44}^{''})^{(9)}((G_{47}),t) =$ **First detrition factor** | |
| $\frac{dG_{13}}{dt} = (a_{13})^{(1)}G_{14} - \begin{bmatrix} (a_{13}^{'})^{(1)}\boxed{+(a_{13}^{''})^{(1)}(T_{14},t)}\boxed{+(a_{16}^{''})^{(2,2,)}(T_{17},t)}\boxed{+(a_{20}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{24}^{''})^{(4,4,4,4,)}(T_{25},t)}\boxed{+(a_{28}^{''})^{(5,5,5,5,)}(T_{29},t)}\boxed{+(a_{32}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{36}^{''})^{(7,7)}(T_{37},t)}\boxed{+(a_{40}^{''})^{(8,8)}(T_{41},t)}\boxed{+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{13}$ | *55* |
| $\frac{dG_{14}}{dt} = (a_{14})^{(1)}G_{13} - \begin{bmatrix} (a_{14}^{'})^{(1)}\boxed{+(a_{14}^{''})^{(1)}(T_{14},t)}\boxed{+(a_{17}^{''})^{(2,2,)}(T_{17},t)}\boxed{+(a_{21}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{25}^{''})^{(4,4,4,4,)}(T_{25},t)}\boxed{+(a_{29}^{''})^{(5,5,5,5,)}(T_{29},t)}\boxed{+(a_{33}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{37}^{''})^{(7,7)}(T_{37},t)}\boxed{+(a_{41}^{''})^{(8,8)}(T_{41},t)}\boxed{+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{14}$ | *56* |
| $\frac{dG_{15}}{dt} = (a_{15})^{(1)}G_{14} - \begin{bmatrix} (a_{15}^{'})^{(1)}\boxed{+(a_{15}^{''})^{(1)}(T_{14},t)}\boxed{+(a_{18}^{''})^{(2,2,)}(T_{17},t)}\boxed{+(a_{22}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{26}^{''})^{(4,4,4,4,)}(T_{25},t)}\boxed{+(a_{30}^{''})^{(5,5,5,5,)}(T_{29},t)}\boxed{+(a_{34}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{38}^{''})^{(7,7)}(T_{37},t)}\boxed{+(a_{42}^{''})^{(8,8)}(T_{41},t)}\boxed{+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{15}$ | *57* |
| Where $\boxed{(a_{13}^{''})^{(1)}(T_{14},t)}$, $\boxed{(a_{14}^{''})^{(1)}(T_{14},t)}$, $\boxed{(a_{15}^{''})^{(1)}(T_{14},t)}$ are first augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{16}^{''})^{(2,2,)}(T_{17},t)}$, $\boxed{+(a_{17}^{''})^{(2,2,)}(T_{17},t)}$, $\boxed{+(a_{18}^{''})^{(2,2,)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{20}^{''})^{(3,3,)}(T_{21},t)}$, $\boxed{+(a_{21}^{''})^{(3,3,)}(T_{21},t)}$, $\boxed{+(a_{22}^{''})^{(3,3,)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{24}^{''})^{(4,4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{25}^{''})^{(4,4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{26}^{''})^{(4,4,4,4,)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{28}^{''})^{(5,5,5,5,)}(T_{29},t)}$, $\boxed{+(a_{29}^{''})^{(5,5,5,5,)}(T_{29},t)}$, $\boxed{+(a_{30}^{''})^{(5,5,5,5,)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{32}^{''})^{(6,6,6,6,)}(T_{33},t)}$, $\boxed{+(a_{33}^{''})^{(6,6,6,6,)}(T_{33},t)}$, $\boxed{+(a_{34}^{''})^{(6,6,6,6,)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{38}^{''})^{(7,7)}(T_{37},t)}$ $\boxed{+(a_{37}^{''})^{(7,7)}(T_{37},t)}$ $\boxed{+(a_{36}^{''})^{(7,7)}(T_{37},t)}$ are seventh augmentation coefficient for 1,2,3<br>$\boxed{+(a_{40}^{''})^{(8,8)}(T_{41},t)}$ $\boxed{+(a_{41}^{''})^{(8,8)}(T_{41},t)}$ $\boxed{+(a_{42}^{''})^{(8,8)}(T_{41},t)}$ are eight augmentation coefficient for 1,2,3<br>$\boxed{+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)}$ are ninth augmentation coefficient for 1,2,3 | |

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| $\frac{dT_{13}}{dt} = (b_{13})^{(1)}T_{14} - \begin{bmatrix} (b_{13}^{'})^{(1)}\boxed{-(b_{13}^{''})^{(1)}(G,t)} & \boxed{-(b_{16}^{''})^{(2,2,)}(G_{19},t)} & \boxed{-(b_{20}^{''})^{(3,3,)}(G_{23},t)} \\ \boxed{-(b_{24}^{''})^{(4,4,4,4,)}(G_{27},t)} & \boxed{-(b_{28}^{''})^{(5,5,5,5,)}(G_{31},t)} & \boxed{-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)} \\ \boxed{-(b_{36}^{''})^{(7,7,)}(G_{39},t)} & \boxed{-(b_{40}^{''})^{(8,8)}(G_{43},t)} & \boxed{-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{13}$ | 58 |
| $\frac{dT_{14}}{dt} = (b_{14})^{(1)}T_{13} - \begin{bmatrix} (b_{14}^{'})^{(1)}\boxed{-(b_{14}^{''})^{(1)}(G,t)} & \boxed{-(b_{17}^{''})^{(2,2,)}(G_{19},t)} & \boxed{-(b_{21}^{''})^{(3,3,)}(G_{23},t)} \\ \boxed{-(b_{25}^{''})^{(4,4,4,4,)}(G_{27},t)} & \boxed{-(b_{29}^{''})^{(5,5,5,5,)}(G_{31},t)} & \boxed{-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)} \\ \boxed{-(b_{37}^{''})^{(7,7,)}(G_{39},t)} & \boxed{-(b_{41}^{''})^{(8,8)}(G_{43},t)} & \boxed{-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{14}$ | 59 |
| $\frac{dT_{15}}{dt} = (b_{15})^{(1)}T_{14} - \begin{bmatrix} (b_{15}^{'})^{(1)}\boxed{-(b_{15}^{''})^{(1)}(G,t)} & \boxed{-(b_{18}^{''})^{(2,2,)}(G_{19},t)} & \boxed{-(b_{22}^{''})^{(3,3,)}(G_{23},t)} \\ \boxed{-(b_{26}^{''})^{(4,4,4,4,)}(G_{27},t)} & \boxed{-(b_{30}^{''})^{(5,5,5,5,)}(G_{31},t)} & \boxed{-(b_{34}^{''})^{(6,6,6,6,)}(G_{35},t)} \\ \boxed{-(b_{38}^{''})^{(7,7,)}(G_{39},t)} & \boxed{-(b_{42}^{''})^{(8,8)}(G_{43},t)} & \boxed{-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{15}$ | 60 |
| Where $\boxed{-(b_{13}^{''})^{(1)}(G,t)}$, $\boxed{-(b_{14}^{''})^{(1)}(G,t)}$, $\boxed{-(b_{15}^{''})^{(1)}(G,t)}$ are first detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{16}^{''})^{(2,2,)}(G_{19},t)}$, $\boxed{-(b_{17}^{''})^{(2,2,)}(G_{19},t)}$, $\boxed{-(b_{18}^{''})^{(2,2,)}(G_{19},t)}$ are second detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{20}^{''})^{(3,3,)}(G_{23},t)}$, $\boxed{-(b_{21}^{''})^{(3,3,)}(G_{23},t)}$, $\boxed{-(b_{22}^{''})^{(3,3,)}(G_{23},t)}$ are third detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{24}^{''})^{(4,4,4,4,)}(G_{27},t)}$, $\boxed{-(b_{25}^{''})^{(4,4,4,4,)}(G_{27},t)}$, $\boxed{-(b_{26}^{''})^{(4,4,4,4,)}(G_{27},t)}$ are fourth detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{28}^{''})^{(5,5,5,5,)}(G_{31},t)}$, $\boxed{-(b_{29}^{''})^{(5,5,5,5,)}(G_{31},t)}$, $\boxed{-(b_{30}^{''})^{(5,5,5,5,)}(G_{31},t)}$ are fifth detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)}$, $\boxed{-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)}$, $\boxed{-(b_{34}^{''})^{(6,6,6,6,)}(G_{35},t)}$ are sixth detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{37}^{''})^{(7,7,)}(G_{39},t)}$, $\boxed{-(b_{36}^{''})^{(7,7,)}(G_{39},t)}$, $\boxed{-(b_{38}^{''})^{(7,7,)}(G_{39},t)}$ are seventh detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{40}^{''})^{(8,8)}(G_{43},t)}$ $\boxed{-(b_{41}^{''})^{(8,8)}(G_{43},t)}$ $\boxed{-(b_{42}^{''})^{(8,8)}(G_{43},t)}$ are eight detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2 and 3 | |
| | |
| $\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \begin{bmatrix} (a_{16}^{'})^{(2)}\boxed{+(a_{16}^{''})^{(2)}(T_{17},t)} & \boxed{+(a_{13}^{''})^{(1,1,)}(T_{14},t)} & \boxed{+(a_{20}^{''})^{(3,3,3)}(T_{21},t)} \\ \boxed{+(a_{24}^{''})^{(4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{28}^{''})^{(5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{32}^{''})^{(6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{36}^{''})^{(7,7,7)}(T_{37},t)} & \boxed{+(a_{40}^{''})^{(8,8,8)}(T_{41},t)} & \boxed{+(a_{44}^{''})^{(9,9)}(T_{45},t)} \end{bmatrix} G_{16}$ | 61 |
| $\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \begin{bmatrix} (a_{17}^{'})^{(2)}\boxed{+(a_{17}^{''})^{(2)}(T_{17},t)} & \boxed{+(a_{14}^{''})^{(1,1,)}(T_{14},t)} & \boxed{+(a_{21}^{''})^{(3,3,3)}(T_{21},t)} \\ \boxed{+(a_{25}^{''})^{(4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{29}^{''})^{(5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{33}^{''})^{(6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{37}^{''})^{(7,7,7)}(T_{37},t)} & \boxed{+(a_{41}^{''})^{(8,8,8)}(T_{41},t)} & \boxed{+(a_{45}^{''})^{(9,9)}(T_{45},t)} \end{bmatrix} G_{17}$ | 62 |

$$\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \begin{bmatrix} (a'_{18})^{(2)}\boxed{+(a''_{18})^{(2)}(T_{17},t)} & \boxed{+(a''_{15})^{(1,1,)}(T_{14},t)} & \boxed{+(a''_{22})^{(3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{26})^{(4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{30})^{(5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{34})^{(6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{38})^{(7,7,7)}(T_{37},t)} & \boxed{+(a''_{42})^{(8,8,8)}(T_{41},t)} & \boxed{+(a''_{46})^{(9,9)}(T_{45},t)} \end{bmatrix} G_{18} \quad (63)$$

Where $\boxed{+(a''_{16})^{(2)}(T_{17},t)}$, $\boxed{+(a''_{17})^{(2)}(T_{17},t)}$, $\boxed{+(a''_{18})^{(2)}(T_{17},t)}$ are first augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{13})^{(1,1,)}(T_{14},t)}$, $\boxed{+(a''_{14})^{(1,1,)}(T_{14},t)}$, $\boxed{+(a''_{15})^{(1,1,)}(T_{14},t)}$ are second augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{20})^{(3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{21})^{(3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{22})^{(3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{24})^{(4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{25})^{(4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{26})^{(4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{28})^{(5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{29})^{(5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{30})^{(5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{32})^{(6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{33})^{(6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{34})^{(6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{36})^{(7,7,7)}(T_{37},t)}$, $\boxed{+(a''_{37})^{(7,7,7)}(T_{37},t)}$, $\boxed{+(a''_{38})^{(7,7,7)}(T_{37},t)}$ are seventh augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{40})^{(8,8,8)}(T_{41},t)}$, $\boxed{+(a''_{41})^{(8,8,8)}(T_{41},t)}$, $\boxed{+(a''_{42})^{(8,8,8)}(T_{41},t)}$ are eight augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{44})^{(9,9)}(T_{45},t)}$, $\boxed{+(a''_{45})^{(9,9)}(T_{45},t)}$, $\boxed{+(a''_{46})^{(9,9)}(T_{45},t)}$ are ninth augmentation coefficient for category 1, 2 and 3

$$\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \begin{bmatrix} (b'_{16})^{(2)}\boxed{-(b''_{16})^{(2)}(G_{19},t)} & \boxed{-(b''_{13})^{(1,1,)}(G,t)} & \boxed{-(b''_{20})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b''_{24})^{(4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{28})^{(5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{32})^{(6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{36})^{(7,7,7)}(G_{39},t)} & \boxed{-(b''_{40})^{(8,8,8)}(G_{43},t)} & \boxed{-(b''_{44})^{(9,9)}(G_{47},t)} \end{bmatrix} T_{16} \quad (64)$$

$$\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \begin{bmatrix} (b'_{17})^{(2)}\boxed{-(b''_{17})^{(2)}(G_{19},t)} & \boxed{-(b''_{14})^{(1,1,)}(G,t)} & \boxed{-(b''_{21})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b''_{25})^{(4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{29})^{(5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{33})^{(6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{37})^{(7,7,7)}(G_{39},t)} & \boxed{-(b''_{41})^{(8,8,8)}(G_{43},t)} & \boxed{-(b''_{45})^{(9,9)}(G_{47},t)} \end{bmatrix} T_{17} \quad (65)$$

$$\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \begin{bmatrix} (b'_{18})^{(2)}\boxed{-(b''_{18})^{(2)}(G_{19},t)} & \boxed{-(b''_{15})^{(1,1,)}(G,t)} & \boxed{-(b''_{22})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b''_{26})^{(4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{38})^{(7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9)}(G_{47},t)} \end{bmatrix} T_{18} \quad (66)$$

where $\boxed{-(b''_{16})^{(2)}(G_{19},t)}$, $\boxed{-(b''_{17})^{(2)}(G_{19},t)}$, $\boxed{-(b''_{18})^{(2)}(G_{19},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{13})^{(1,1,)}(G,t)}$, $\boxed{-(b''_{14})^{(1,1,)}(G,t)}$, $\boxed{-(b''_{15})^{(1,1,)}(G,t)}$ are second detrition coefficients for category 1,2 and 3

$\boxed{-(b''_{20})^{(3,3,3,)}(G_{23},t)}$, $\boxed{-(b''_{21})^{(3,3,3,)}(G_{23},t)}$, $\boxed{-(b''_{22})^{(3,3,3,)}(G_{23},t)}$ are third detrition coefficients for category 1,2 and 3

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| $-(b_{24}^{''})^{(4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1,2 and 3<br>$-(b_{28}^{''})^{(5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1,2 and 3<br>$-(b_{32}^{''})^{(6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6)}(G_{35},t)$ are sixth detrition coefficients for category 1,2 and 3<br>$-(b_{36}^{''})^{(7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1,2 and 3<br>$-(b_{40}^{''})^{(8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8)}(G_{43},t)$ are eight detrition coefficients for category 1,2 and 3<br>$-(b_{44}^{''})^{(9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1,2 and 3 | |
| $\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \left[ \begin{array}{c} (a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21},t)\ +(a_{16}^{''})^{(2,2,2)}(T_{17},t)\ +(a_{13}^{''})^{(1,1,1,)}(T_{14},t) \\ +(a_{24}^{''})^{(4,4,4,4,4,4)}(T_{25},t)\ +(a_{28}^{''})^{(5,5,5,5,5,5)}(T_{29},t)\ +(a_{32}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ +(a_{36}^{''})^{(7,7,7,7)}(T_{37},t)\ +(a_{40}^{''})^{(8,8,8,8)}(T_{41},t)\ +(a_{44}^{''})^{(9,9,9)}(T_{45},t) \end{array} \right] G_{20}$ | 67 |
| $\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \left[ \begin{array}{c} (a_{21}^{'})^{(3)}+(a_{21}^{''})^{(3)}(T_{21},t)\ +(a_{17}^{''})^{(2,2,2)}(T_{17},t)\ +(a_{14}^{''})^{(1,1,1,)}(T_{14},t) \\ +(a_{25}^{''})^{(4,4,4,4,4,4)}(T_{25},t)\ +(a_{29}^{''})^{(5,5,5,5,5,5)}(T_{29},t)\ +(a_{33}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ +(a_{37}^{''})^{(7,7,7,7)}(T_{37},t)\ +(a_{41}^{''})^{(8,8,8,8)}(T_{41},t)\ +(a_{45}^{''})^{(9,9,9)}(T_{45},t) \end{array} \right] G_{21}$ | 68 |
| $\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \left[ \begin{array}{c} (a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21},t)\ +(a_{18}^{''})^{(2,2,2)}(T_{17},t)\ +(a_{15}^{''})^{(1,1,1,)}(T_{14},t) \\ +(a_{26}^{''})^{(4,4,4,4,4,4)}(T_{25},t)\ +(a_{30}^{''})^{(5,5,5,5,5,5)}(T_{29},t)\ +(a_{34}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ +(a_{38}^{''})^{(7,7,7,7)}(T_{37},t)\ +(a_{42}^{''})^{(8,8,8,8)}(T_{41},t)\ +(a_{46}^{''})^{(9,9,9)}(T_{45},t) \end{array} \right] G_{22}$ | 69 |
| $+(a_{20}^{''})^{(3)}(T_{21},t)$, $+(a_{21}^{''})^{(3)}(T_{21},t)$, $+(a_{22}^{''})^{(3)}(T_{21},t)$ are first augmentation coefficients for category 1, 2 and 3<br>$+(a_{16}^{''})^{(2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2)}(T_{17},t)$ are second augmentation coefficients for category 1, 2 and 3<br>$+(a_{13}^{''})^{(1,1,1,)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,)}(T_{14},t)$ are third augmentation coefficients for category 1, 2 and 3<br>$+(a_{24}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficients for category 1, 2 and 3<br>$+(a_{28}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficients for category 1, 2 and 3<br>$+(a_{32}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation coefficients for category 1, 2 and 3<br>$+(a_{36}^{''})^{(7,7,7,7)}(T_{37},t)$, $+(a_{37}^{''})^{(7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7)}(T_{37},t)$ are seventh augmentation coefficients for category 1, 2 and 3<br>$+(a_{40}^{''})^{(8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8)}(T_{41},t)$ are eight augmentation coefficients | |

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| for category 1, 2 and 3<br>$+(a_{44}^{''})^{(9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9)}(T_{45},t)$, $+(a_{46}^{''})^{(9,9,9)}(T_{45},t)$ are ninth augmentation coefficients for category 1, 2 and 3 | |
| $\frac{dT_{20}}{dt} = (b_{20})^{(3)}T_{21} - \begin{bmatrix} (b_{20}^{'})^{(3)}\,-(b_{20}^{''})^{(3)}(G_{23},t) & -(b_{16}^{''})^{(2,2,2)}(G_{19},t) & -(b_{13}^{''})^{(1,1,1,)}(G,t) \\ -(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t) & -(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t) & -(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t) \\ -(b_{36}^{''})^{(7,7,7,7)}(G_{39},t) & -(b_{40}^{''})^{(8,8,8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9,9)}(G_{47},t) \end{bmatrix} T_{20}$ | *70* |
| $\frac{dT_{21}}{dt} = (b_{21})^{(3)}T_{20} - \begin{bmatrix} (b_{21}^{'})^{(3)}\,-(b_{21}^{''})^{(3)}(G_{23},t) & -(b_{17}^{''})^{(2,2,2)}(G_{19},t) & -(b_{14}^{''})^{(1,1,1,)}(G,t) \\ -(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t) & -(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t) & -(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t) \\ -(b_{37}^{''})^{(7,7,7,7)}(G_{39},t) & -(b_{41}^{''})^{(8,8,8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9,9)}(G_{47},t) \end{bmatrix} T_{21}$ | *71* |
| $\frac{dT_{22}}{dt} = (b_{22})^{(3)}T_{21} - \begin{bmatrix} (b_{22}^{'})^{(3)}\,-(b_{22}^{''})^{(3)}(G_{23},t) & -(b_{18}^{''})^{(2,2,2)}(G_{19},t) & -(b_{15}^{''})^{(1,1,1,)}(G,t) \\ -(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t) & -(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t) & -(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t) \\ -(b_{38}^{''})^{(7,7,7,7)}(G_{39},t) & -(b_{42}^{''})^{(8,8,8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9)}(G_{47},t) \end{bmatrix} T_{22}$ | *72* |
| $-(b_{20}^{''})^{(3)}(G_{23},t)$, $-(b_{21}^{''})^{(3)}(G_{23},t)$, $-(b_{22}^{''})^{(3)}(G_{23},t)$ *are first detrition coefficients for category 1, 2 and 3*<br>$-(b_{16}^{''})^{(2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2)}(G_{19},t)$ *are second detrition coefficients for category 1, 2 and 3*<br>$-(b_{13}^{''})^{(1,1,1,)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,)}(G,t)$, $-(b_{15}^{''})^{(1,1,1,)}(G,t)$ *are third detrition coefficients for category 1,2 and 3*<br>$-(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$ *are fourth detrition coefficients for category 1, 2 and 3*<br>$-(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$ *are fifth detrition coefficients for category 1, 2 and 3*<br>$-(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$ *are sixth detrition coefficients for category 1, 2 and 3*<br>$-(b_{36}^{''})^{(7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7)}(G_{39},t)$ $-(b_{38}^{''})^{(7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3<br>$-(b_{40}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8)}(G_{43},t)$ are eight detrition coefficients for category 1, 2 and 3<br>$-(b_{46}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3 | |
| $\frac{dG_{24}}{dt} = (a_{24})^{(4)}G_{25} - \begin{bmatrix} (a_{24}^{'})^{(4)}\,+(a_{24}^{''})^{(4)}(T_{25},t) & +(a_{28}^{''})^{(5,5,)}(T_{29},t) & +(a_{32}^{''})^{(6,6,)}(T_{33},t) \\ +(a_{13}^{''})^{(1,1,1,1)}(T_{14},t) & +(a_{16}^{''})^{(2,2,2,2)}(T_{17},t) & +(a_{20}^{''})^{(3,3,3,3)}(T_{21},t) \\ +(a_{36}^{''})^{(7,7,7,7,7)}(T_{37},t) & +(a_{40}^{''})^{(8,8,8,8,8)}(T_{41},t) & +(a_{44}^{''})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{24}$ | *73* |

$$\frac{dG_{25}}{dt} = (a_{25})^{(4)}G_{24} - \begin{bmatrix} (a'_{25})^{(4)} & +(a''_{25})^{(4)}(T_{25},t) & +(a''_{29})^{(5,5,)}(T_{29},t) & +(a''_{33})^{(6,6)}(T_{33},t) \\ +(a''_{14})^{(1,1,1,1)}(T_{14},t) & +(a''_{17})^{(2,2,2,2)}(T_{17},t) & +(a''_{21})^{(3,3,3,3)}(T_{21},t) \\ +(a''_{37})^{(7,7,7,7,7)}(T_{37},t) & +(a''_{41})^{(8,8,8,8,8)}(T_{41},t) & +(a''_{45})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{25} \quad (74)$$

$$\frac{dG_{26}}{dt} = (a_{26})^{(4)}G_{25} - \begin{bmatrix} (a'_{26})^{(4)} & +(a''_{26})^{(4)}(T_{25},t) & +(a''_{30})^{(5,5,)}(T_{29},t) & +(a''_{34})^{(6,6,)}(T_{33},t) \\ +(a''_{15})^{(1,1,1,1)}(T_{14},t) & +(a''_{18})^{(2,2,2,2)}(T_{17},t) & +(a''_{22})^{(3,3,3,3)}(T_{21},t) \\ +(a''_{38})^{(7,7,7,7,7)}(T_{37},t) & +(a''_{42})^{(8,8,8,8,8)}(T_{41},t) & +(a''_{46})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{26} \quad (75)$$

$(a''_{24})^{(4)}(T_{25},t)$, $(a''_{25})^{(4)}(T_{25},t)$, $(a''_{26})^{(4)}(T_{25},t)$ *are first augmentation coefficients category* 1, 2 3

$+(a''_{28})^{(5,5,)}(T_{29},t)$, $+(a''_{29})^{(5,5,)}(T_{29},t)$, $+(a''_{30})^{(5,5,)}(T_{29},t)$ *are second augmentation coefficient for category* 1, 2 *and* 3

$+(a''_{32})^{(6,6,)}(T_{33},t)$, $+(a''_{33})^{(6,6,)}(T_{33},t)$, $+(a''_{34})^{(6,6,)}(T_{33},t)$ *are third augmentation coefficient for category* 1, 2 *and* 3

$+(a''_{13})^{(1,1,1,1)}(T_{14},t)$, $+(a''_{14})^{(1,1,1,1)}(T_{14},t)$, $+(a''_{15})^{(1,1,1,1)}(T_{14},t)$ *are fourth augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{16})^{(2,2,2,2)}(T_{17},t)$, $+(a''_{17})^{(2,2,2,2)}(T_{17},t)$, $+(a''_{18})^{(2,2,2,2)}(T_{17},t)$ *are fifth augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{20})^{(3,3,3,3)}(T_{21},t)$, $+(a''_{21})^{(3,3,3,3)}(T_{21},t)$, $+(a''_{22})^{(3,3,3,3)}(T_{21},t)$ *are sixth augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{36})^{(7,7,7,7,7)}(T_{37},t)$, $+(a''_{37})^{(7,7,7,7,7)}(T_{37},t)$, $+(a''_{38})^{(7,7,7,7,7)}(T_{37},t)$ *are seventh augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{40})^{(8,8,8,8,8)}(T_{41},t)$, $+(a''_{41})^{(8,8,8,8,8)}(T_{41},t)$, $+(a''_{42})^{(8,8,8,8,8)}(T_{41},t)$ *are eighth augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{46})^{(9,9,9,9)}(T_{45},t)$, $+(a''_{45})^{(9,9,9,9)}(T_{45},t)$, $+(a''_{44})^{(9,9,9,9)}(T_{45},t)$ are ninth detrition coefficients for category 1 2 3

$$\frac{dT_{24}}{dt} = (b_{24})^{(4)}T_{25} - \begin{bmatrix} (b'_{24})^{(4)} & -(b''_{24})^{(4)}(G_{27},t) & -(b''_{28})^{(5,5,)}(G_{31},t) & -(b''_{32})^{(6,6,)}(G_{35},t) \\ -(b''_{13})^{(1,1,1,1)}(G,t) & -(b''_{16})^{(2,2,2,2)}(G_{19},t) & -(b''_{20})^{(3,3,3,3)}(G_{23},t) \\ -(b''_{36})^{(7,7,7,7,7)}(G_{39},t) & -(b''_{40})^{(8,8,8,8,8)}(G_{43},t) & -(b''_{44})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{24} \quad (76)$$

$$\frac{dT_{25}}{dt} = (b_{25})^{(4)}T_{24} - \begin{bmatrix} (b'_{25})^{(4)} & -(b''_{25})^{(4)}(G_{27},t) & -(b''_{29})^{(5,5,)}(G_{31},t) & -(b''_{33})^{(6,6,)}(G_{35},t) \\ -(b''_{14})^{(1,1,1,1)}(G,t) & -(b''_{17})^{(2,2,2,2)}(G_{19},t) & -(b''_{21})^{(3,3,3,3)}(G_{23},t) \\ -(b''_{37})^{(7,7,7,7,7)}(G_{39},t) & -(b''_{41})^{(8,8,8,8,8)}(G_{43},t) & -(b''_{45})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{25} \quad (77)$$

$$\frac{dT_{26}}{dt} = (b_{26})^{(4)}T_{25} - \begin{bmatrix} (b'_{26})^{(4)} & -(b''_{26})^{(4)}(G_{27},t) & -(b''_{30})^{(5,5,)}(G_{31},t) & -(b''_{34})^{(6,6,)}(G_{35},t) \\ -(b''_{15})^{(1,1,1,1)}(G,t) & -(b''_{18})^{(2,2,2,2)}(G_{19},t) & -(b''_{22})^{(3,3,3,3)}(G_{23},t) \\ -(b''_{38})^{(7,7,7,7,7)}(G_{39},t) & -(b''_{42})^{(8,8,8,8,8)}(G_{43},t) & -(b''_{46})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{26} \quad (78)$$

*Where* $-(b''_{24})^{(4)}(G_{27},t)$, $-(b''_{25})^{(4)}(G_{27},t)$, $-(b''_{26})^{(4)}(G_{27},t)$ *are first detrition coefficients for category* 1, 2 *and* 3

$-(b_{28}^{''})^{(5,5,)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,)}(G_{31},t)$ *are second detrition coefficients for category* 1, 2 *and* 3

$-(b_{32}^{''})^{(6,6,)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,)}(G_{35},t)$ *are third detrition coefficients for category* 1, 2 *and* 3

$-(b_{13}^{''})^{(1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1)}(G,t)$, $-(b_{15}^{''})^{(1,1,1,1)}(G,t)$ *are fourth detrition coefficients for category* 1, 2 *and* 3

$-(b_{16}^{''})^{(2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2)}(G_{19},t)$ *are fifth detrition coefficients for category* 1, 2 *and* 3

$-(b_{20}^{''})^{(3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3)}(G_{23},t)$ *are sixth detrition coefficients for category* 1, 2 *and* 3

$-(b_{36}^{''})^{(7,7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7,7)}(G_{39},t)$ *are seventh detrition coefficients for category* 1, 2 *and* 3

$-(b_{40}^{''})^{(8,8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8,8)}(G_{43},t)$ *are eighth detrition coefficients for category* 1, 2 *and* 3

$-(b_{46}^{''})^{(9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1 2 3

$$\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \left[\begin{array}{ccc} (a_{28}^{'})^{(5)} & +(a_{28}^{''})^{(5)}(T_{29},t) & +(a_{24}^{''})^{(4,4,)}(T_{25},t) \quad +(a_{32}^{''})^{(6,6,6)}(T_{33},t) \\ +(a_{13}^{''})^{(1,1,1,1,1)}(T_{14},t) & +(a_{16}^{''})^{(2,2,2,2,2)}(T_{17},t) & +(a_{20}^{''})^{(3,3,3,3,3)}(T_{21},t) \\ +(a_{36}^{''})^{(7,7,7,7,7)}(T_{37},t) & +(a_{40}^{''})^{(8,8,8,8,8)}(T_{41},t) & +(a_{44}^{''})^{(9,9,9,9,9)}(T_{45},t) \end{array}\right] G_{28} \quad (79)$$

$$\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \left[\begin{array}{ccc} (a_{29}^{'})^{(5)} & +(a_{29}^{''})^{(5)}(T_{29},t) & +(a_{25}^{''})^{(4,4,)}(T_{25},t) \quad +(a_{33}^{''})^{(6,6,6)}(T_{33},t) \\ +(a_{14}^{''})^{(1,1,1,1,1)}(T_{14},t) & +(a_{17}^{''})^{(2,2,2,2,2)}(T_{17},t) & +(a_{21}^{''})^{(3,3,3,3,3)}(T_{21},t) \\ +(a_{37}^{''})^{(7,7,7,7,7)}(T_{37},t) & +(a_{41}^{''})^{(8,8,8,8,8)}(T_{41},t) & +(a_{45}^{''})^{(9,9,9,9,9)}(T_{45},t) \end{array}\right] G_{29} \quad (80)$$

$$\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \left[\begin{array}{ccc} (a_{30}^{'})^{(5)} & +(a_{30}^{''})^{(5)}(T_{29},t) & +(a_{26}^{''})^{(4,4,)}(T_{25},t) \quad +(a_{34}^{''})^{(6,6,6)}(T_{33},t) \\ +(a_{15}^{''})^{(1,1,1,1,1)}(T_{14},t) & +(a_{18}^{''})^{(2,2,2,2,2)}(T_{17},t) & +(a_{22}^{''})^{(3,3,3,3,3)}(T_{21},t) \\ +(a_{38}^{''})^{(7,7,7,7,7)}(T_{37},t) & +(a_{42}^{''})^{(8,8,8,8,8)}(T_{41},t) & +(a_{46}^{''})^{(9,9,9,9,9)}(T_{45},t) \end{array}\right] G_{30} \quad (81)$$

*Where* $+(a_{28}^{''})^{(5)}(T_{29},t)$, $+(a_{29}^{''})^{(5)}(T_{29},t)$, $+(a_{30}^{''})^{(5)}(T_{29},t)$ *are first augmentation coefficients for category* 1, 2 *and* 3

*And* $+(a_{24}^{''})^{(4,4,)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,)}(T_{25},t)$ *are second augmentation coefficient for category* 1, 2 *and* 3

$+(a_{32}^{''})^{(6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6)}(T_{33},t)$ *are third augmentation coefficient for category* 1, 2 *and* 3

$+(a_{13}^{''})^{(1,1,1,1,1)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1,1)}(T_{14},t)$ are fourth augmentation coefficients for category 1,2, and 3

$+(a_{16}^{''})^{(2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2)}(T_{17},t)$ are fifth augmentation coefficients for category 1,2,and 3

$+(a_{20}^{''})^{(3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3)}(T_{21},t)$ are sixth augmentation

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| coefficients for category 1,2, 3<br>$+(a_{36}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)$ , $+(a_{37}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)$ are seventh augmentation coefficients for category 1,2, 3<br>$+(a_{40}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)$ are eighth augmentation coefficients for category 1,2, 3<br>$+(a_{46}^{''})^{(9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficients for category 1,2, 3 | |
| $\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \begin{bmatrix} (b_{28}^{'})^{(5)} & -(b_{28}^{''})^{(5)}(G_{31},t) & -(b_{24}^{''})^{(4,4,)}(G_{27},t) & -(b_{32}^{''})^{(6,6,6)}(G_{35},t) \\ & -(b_{13}^{''})^{(1,1,1,1,1)}(G,t) & -(b_{16}^{''})^{(2,2,2,2,2)}(G_{19},t) & -(b_{20}^{''})^{(3,3,3,3,3)}(G_{23},t) \\ & -(b_{36}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b_{40}^{''})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{28}$ | *82* |
| $\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \begin{bmatrix} (b_{29}^{'})^{(5)} & -(b_{29}^{''})^{(5)}(G_{31},t) & -(b_{25}^{''})^{(4,4,)}(G_{27},t) & -(b_{33}^{''})^{(6,6,6)}(G_{35},t) \\ & -(b_{14}^{''})^{(1,1,1,1,1)}(G,t) & -(b_{17}^{''})^{(2,2,2,2,2)}(G_{19},t) & -(b_{21}^{''})^{(3,3,3,3,3)}(G_{23},t) \\ & -(b_{37}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b_{41}^{''})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{29}$ | *83* |
| $\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \begin{bmatrix} (b_{30}^{'})^{(5)} & -(b_{30}^{''})^{(5)}(G_{31},t) & -(b_{26}^{''})^{(4,4,)}(G_{27},t) & -(b_{34}^{''})^{(6,6,6)}(G_{35},t) \\ & -(b_{15}^{''})^{(1,1,1,1,1,)}(G,t) & -(b_{18}^{''})^{(2,2,2,2,2)}(G_{19},t) & -(b_{22}^{''})^{(3,3,3,3,3)}(G_{23},t) \\ & -(b_{38}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b_{42}^{''})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{30}$ | *84* |
| *where* $-(b_{28}^{''})^{(5)}(G_{31},t)$ , $-(b_{29}^{''})^{(5)}(G_{31},t)$ , $-(b_{30}^{''})^{(5)}(G_{31},t)$ *are first detrition coefficients for category* 1, 2 *and* 3<br>$-(b_{24}^{''})^{(4,4,)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,)}(G_{27},t)$ *are second detrition coefficients for category* 1,2 *and* 3<br>$-(b_{32}^{''})^{(6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6)}(G_{35},t)$ *are third detrition coefficients for category* 1,2 *and* 3<br>$-(b_{13}^{''})^{(1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1)}(G,t)$ , $-(b_{15}^{''})^{(1,1,1,1,1,)}(G,t)$ are fourth detrition coefficients for category 1,2, and 3<br>$-(b_{16}^{''})^{(2,2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2,2)}(G_{19},t)$ are fifth detrition coefficients for category 1,2, and 3<br>$-(b_{20}^{''})^{(3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3)}(G_{23},t)$ are sixth detrition coefficients for category 1,2, and 3<br>$-(b_{36}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1,2, and 3<br>$-(b_{42}^{''})^{(8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{40}^{''})^{(8,8,8,8,8,8,8)}(G_{43},t)$ are eighth detrition coefficients for category 1,2, and 3<br>$-(b_{46}^{''})^{(9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1,2, and 3 | |

$$\frac{dG_{32}}{dt} = (a_{32})^{(6)} G_{33} - \begin{bmatrix} (a'_{32})^{(6)} & +(a''_{32})^{(6)}(T_{33},t) & +(a''_{28})^{(5,5,5)}(T_{29},t) & +(a''_{24})^{(4,4,4,)}(T_{25},t) \\ & +(a''_{13})^{(1,1,1,1,1,1)}(T_{14},t) & +(a''_{16})^{(2,2,2,2,2,2)}(T_{17},t) & +(a''_{20})^{(3,3,3,3,3,3)}(T_{21},t) \\ & +(a''_{36})^{(7,7,7,7,7,7,7)}(T_{37},t) & +(a''_{40})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a''_{44})^{(9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{32} \quad (85)$$

$$\frac{dG_{33}}{dt} = (a_{33})^{(6)} G_{32} - \begin{bmatrix} (a'_{33})^{(6)} & +(a''_{33})^{(6)}(T_{33},t) & +(a''_{29})^{(5,5,5)}(T_{29},t) & +(a''_{25})^{(4,4,4,)}(T_{25},t) \\ & +(a''_{14})^{(1,1,1,1,1,1)}(T_{14},t) & +(a''_{17})^{(2,2,2,2,2,2)}(T_{17},t) & +(a''_{21})^{(3,3,3,3,3,3)}(T_{21},t) \\ & +(a''_{37})^{(7,7,7,7,7,7,7)}(T_{37},t) & +(a''_{41})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a''_{45})^{(9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{33} \quad (86)$$

$$\frac{dG_{34}}{dt} = (a_{34})^{(6)} G_{33} - \begin{bmatrix} (a'_{34})^{(6)} & +(a''_{34})^{(6)}(T_{33},t) & +(a''_{30})^{(5,5,5)}(T_{29},t) & +(a''_{26})^{(4,4,4,)}(T_{25},t) \\ & +(a''_{15})^{(1,1,1,1,1,1)}(T_{14},t) & +(a''_{18})^{(2,2,2,2,2,2)}(T_{17},t) & +(a''_{22})^{(3,3,3,3,3,3)}(T_{21},t) \\ & +(a''_{38})^{(7,7,7,7,7,7,7)}(T_{37},t) & +(a''_{42})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a''_{46})^{(9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{34} \quad (87)$$

$+(a''_{32})^{(6)}(T_{33},t)$, $+(a''_{33})^{(6)}(T_{33},t)$, $+(a''_{34})^{(6)}(T_{33},t)$ *are first augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{28})^{(5,5,5)}(T_{29},t)$, $+(a''_{29})^{(5,5,5)}(T_{29},t)$, $+(a''_{30})^{(5,5,5)}(T_{29},t)$ *are second augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{24})^{(4,4,4,)}(T_{25},t)$, $+(a''_{25})^{(4,4,4,)}(T_{25},t)$, $+(a''_{26})^{(4,4,4,)}(T_{25},t)$ *are third augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{13})^{(1,1,1,1,1,1)}(T_{14},t)$, $+(a''_{14})^{(1,1,1,1,1,1)}(T_{14},t)$, $+(a''_{15})^{(1,1,1,1,1,1)}(T_{14},t)$ - are fourth augmentation coefficients

$+(a''_{16})^{(2,2,2,2,2,2)}(T_{17},t)$, $+(a''_{17})^{(2,2,2,2,2,2)}(T_{17},t)$, $+(a''_{18})^{(2,2,2,2,2,2)}(T_{17},t)$ - fifth augmentation coefficients

$+(a''_{20})^{(3,3,3,3,3,3)}(T_{21},t)$, $+(a''_{21})^{(3,3,3,3,3,3)}(T_{21},t)$, $+(a''_{22})^{(3,3,3,3,3,3)}(T_{21},t)$ sixth augmentation coefficients

$+(a''_{36})^{(7,7,7,7,7,7,7)}(T_{37},t)$, $+(a''_{37})^{(7,7,7,7,7,7,7)}(T_{37},t)$, $+(a''_{38})^{(7,7,7,7,7,7,7)}(T_{37},t)$ seventh augmentation coefficients

$+(a''_{40})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a''_{41})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a''_{42})^{(8,8,8,8,8,8,8)}(T_{41},t)$ Eighth augmentation coefficients

$+(a''_{44})^{(9,9,9,9,9,9)}(T_{45},t)$, $+(a''_{45})^{(9,9,9,9,9,9)}(T_{45},t)$, $+(a''_{46})^{(9,9,9,9,9,9)}(T_{45},t)$ ninth augmentation coefficients

$$\frac{dT_{32}}{dt} = (b_{32})^{(6)} T_{33} - \begin{bmatrix} (b'_{32})^{(6)} & -(b''_{32})^{(6)}(G_{35},t) & -(b''_{28})^{(5,5,5)}(G_{31},t) & -(b''_{24})^{(4,4,4,)}(G_{27},t) \\ & -(b''_{13})^{(1,1,1,1,1,1)}(G,t) & -(b''_{16})^{(2,2,2,2,2,2)}(G_{19},t) & -(b''_{20})^{(3,3,3,3,3,3)}(G_{23},t) \\ & -(b''_{36})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b''_{40})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b''_{44})^{(9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{32} \quad (88)$$

$$\frac{dT_{33}}{dt} = (b_{33})^{(6)} T_{32} - \begin{bmatrix} (b'_{33})^{(6)} & -(b''_{33})^{(6)}(G_{35},t) & -(b''_{29})^{(5,5,5)}(G_{31},t) & -(b''_{25})^{(4,4,4,)}(G_{27},t) \\ & -(b''_{14})^{(1,1,1,1,1,1)}(G,t) & -(b''_{17})^{(2,2,2,2,2,2)}(G_{19},t) & -(b''_{21})^{(3,3,3,3,3,3)}(G_{23},t) \\ & -(b''_{37})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b''_{41})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b''_{45})^{(9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{33} \quad (89)$$

$$\frac{dT_{34}}{dt} = (b_{34})^{(6)}T_{33} - \begin{bmatrix} (b'_{34})^{(6)}\boxed{-(b''_{34})^{(6)}(G_{35},t)} & \boxed{-(b''_{30})^{(5,5,5)}(G_{31},t)} & \boxed{-(b''_{26})^{(4,4,4,)}(G_{27},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{18})^{(2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{38})^{(7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{34} \quad (90)$$

$\boxed{-(b''_{32})^{(6)}(G_{35},t)}$, $\boxed{-(b''_{33})^{(6)}(G_{35},t)}$, $\boxed{-(b''_{34})^{(6)}(G_{35},t)}$ *are first detrition coefficients for category* 1, 2 *and* 3

$\boxed{-(b''_{28})^{(5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{29})^{(5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{30})^{(5,5,5)}(G_{31},t)}$ *are second detrition coefficients for category* 1, 2 *and* 3

$\boxed{-(b''_{24})^{(4,4,4,)}(G_{27},t)}$, $\boxed{-(b''_{25})^{(4,4,4,)}(G_{27},t)}$, $\boxed{-(b''_{26})^{(4,4,4,)}(G_{27},t)}$ *are third detrition coefficients for category* 1,2 *and* 3

$\boxed{-(b''_{13})^{(1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{14})^{(1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b''_{15})^{(1,1,1,1,1,1)}(G,t)}$ are fourth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{16})^{(2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{17})^{(2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{18})^{(2,2,2,2,2,2)}(G_{19},t)}$ are fifth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{20})^{(3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{21})^{(3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b''_{22})^{(3,3,3,3,3,3)}(G_{23},t)}$ are sixth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{36})^{(7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{37})^{(7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{38})^{(7,7,7,7,7,7,7)}(G_{39},t)}$ are seventh detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{40})^{(8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{41})^{(8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{42})^{(8,8,8,8,8,8,8,8)}(G_{43},t)}$ are eighth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{46})^{(9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{45})^{(9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{44})^{(9,9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2, and 3

$$\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \begin{bmatrix} (a'_{36})^{(7)}\boxed{+(a''_{36})^{(7)}(T_{37},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{24})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8,)}(T_{41},t)} & \boxed{+(a''_{44})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{13} \quad (91)$$

$$\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \begin{bmatrix} (a'_{37})^{(7)}\boxed{+(a''_{37})^{(7)}(T_{37},t)} & \boxed{+(a''_{17})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{21})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{25})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{29})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{33})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{41})^{(8,8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{45})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{14} \quad (92)$$

$$\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \begin{bmatrix} (a'_{38})^{(7)}\boxed{+(a''_{38})^{(7)}(T_{37},t)} & \boxed{+(a''_{18})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{22})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{26})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{30})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{34})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{15})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{42})^{(8,8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{46})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{15} \quad (93)$$

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| Where $(a_{36}^{''})^{(7)}(T_{37},t)$, $(a_{37}^{''})^{(7)}(T_{37},t)$, $(a_{38}^{''})^{(7)}(T_{37},t)$ are first augmentation coefficients for category 1, 2 and 3<br>$+(a_{16}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t)$ are second augmentation coefficient for category 1, 2 and 3<br>$+(a_{20}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3<br>$+(a_{24}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3<br>$+(a_{28}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficient for category 1, 2 and 3<br>$+(a_{32}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation coefficient for category 1, 2 and 3<br>$+(a_{13}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{13}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t)$ are seventh augmentation coefficient for category 1, 2 and 3<br>$+(a_{42}^{''})^{(8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{40}^{''})^{(8,8,8,8,8,8,8,8,)}(T_{41},t)$ are eighth augmentation coefficient for 1,2,3<br>$+(a_{46}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficient for 1,2,3 | |
| $\frac{dT_{36}}{dt} =$<br>$(b_{36})^{(7)}T_{37} - \begin{bmatrix} (b_{36}^{'})^{(7)}\boxed{-(b_{36}^{''})^{(7)}(G_{39},t)} & \boxed{-(b_{16}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{20}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{24}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{28}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{32}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{13}^{''})^{(1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{40}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b_{44}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{13}$ | 94 |
| $\frac{dT_{37}}{dt} =$<br>$(b_{37})^{(7)}T_{36} - \begin{bmatrix} (b_{37}^{'})^{(7)}\boxed{-(b_{37}^{''})^{(7)}(G_{39},t)} & \boxed{-(b_{17}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{21}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{25}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{29}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{33}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{14}^{''})^{(1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{41}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b_{45}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{14}$ | |
| $\frac{dT_{38}}{dt} =$<br>$(b_{38})^{(7)}T_{37} - \begin{bmatrix} (b_{38}^{'})^{(7)}\boxed{-(b_{38}^{''})^{(7)}(G_{39},t)} & \boxed{-(b_{18}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{22}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{26}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{30}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{34}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{15}^{''})^{(1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{42}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b_{46}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{15}$ | |
| Where $\boxed{-(b_{36}^{''})^{(7)}(G_{39},t)}$, $\boxed{-(b_{37}^{''})^{(7)}(G_{39},t)}$, $\boxed{-(b_{38}^{''})^{(7)}(G_{39},t)}$ are first detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{16}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b_{17}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b_{18}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)}$ are second detrition coefficients for category 1, 2 and 3 | |

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| $-(b_{20}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3<br>$-(b_{24}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3<br>$-(b_{28}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3<br>$-(b_{32}^{''})^{(6,6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6,6,6)}(G_{35},t)$ are sixth detrition coefficients for category 1, 2 and 3<br>$-(b_{15}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{13}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$ are seventh detrition coefficients for category 1, 2 and 3<br>$-(b_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)$ are eighth detrition coefficients for category 1, 2 and 3<br>$-(b_{46}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3 | |
| $\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \begin{bmatrix} (a_{40}^{'})^{(8)}+(a_{40}^{''})^{(8)}(T_{41},t) & +(a_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t) & +(a_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t) & +(a_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t) & +(a_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t) & +(a_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t) & +(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{13}$ | *95* |
| $\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \begin{bmatrix} (a_{41}^{'})^{(8)}+(a_{41}^{''})^{(8)}(T_{41},t) & +(a_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t) & +(a_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t) & +(a_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t) & +(a_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t) & +(a_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t) & +(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{14}$ | |
| $\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \begin{bmatrix} (a_{42}^{'})^{(8)}+(a_{42}^{''})^{(8)}(T_{41},t) & +(a_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t) & +(a_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t) & +(a_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t) & +(a_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t) & +(a_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t) & +(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{15}$ | |
| Where $+(a_{40}^{''})^{(8)}(T_{41},t)$, $+(a_{41}^{''})^{(8)}(T_{41},t)$, $+(a_{42}^{''})^{(8)}(T_{41},t)$ are first augmentation coefficients for category 1, 2 and 3<br>$+(a_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)$ are second augmentation coefficient for category 1, 2 and 3<br>$+(a_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3 | |

$+(a_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$ , $+(a_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$ , $+(a_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3

$+(a_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$ , $+(a_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficient for category 1, 2 and 3

$+(a_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$ , $+(a_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation coefficient for category 1, 2 and 3

$+(a_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)$ $+(a_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)$ $+(a_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)$ are seventh augmentation coefficient for 1,2,3

$+(a_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)$ ,$+(a_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)$ ,$+(a_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)$ are eighth augmentation coefficient for 1,2,3

$+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ ,$+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ ,$+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{40}}{dt} =$$

$$(b_{40})^{(8)}T_{41} - \begin{bmatrix} (b_{40}^{'})^{(8)}\left[-(b_{40}^{''})^{(8)}(G_{43},t)\right] & -(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t) & -(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t) \\ -(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t) & -(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t) & -(b_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t) \\ -(b_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t) & -(b_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t) & -(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{13}$$

$$\frac{dT_{41}}{dt} =$$

$$(b_{41})^{(8)}T_{40} - \begin{bmatrix} (b_{41}^{'})^{(8)}\left[-(b_{41}^{''})^{(8)}(G_{43},t)\right] & -(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t) & -(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t) \\ -(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t) & -(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t) & -(b_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t) \\ -(b_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t) & -(b_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t) & -(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{14}$$

$$\frac{dT_{42}}{dt} =$$

$$(b_{42})^{(8)}T_{41} - \begin{bmatrix} (b_{42}^{'})^{(8)}\left[-(b_{42}^{''})^{(8)}(G_{43},t)\right] & -(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t) & -(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t) \\ -(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t) & -(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t) & -(b_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t) \\ -(b_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t) & -(b_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t) & -(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{15}$$

Where $-(b_{36}^{''})^{(7)}(G_{39},t)$ , $-(b_{37}^{''})^{(7)}(G_{39},t)$ , $-(b_{38}^{''})^{(7)}(G_{39},t)$ are first detrition coefficients for category 1, 2 and 3

$-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)$ , $-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)$ , $-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3

$-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)$ , $-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)$ , $-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3

$-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)$ , $-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)$ , $-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3

$-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)$ , $-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)$ , $-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3

$-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{15}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$ are sixth detrition coefficients for category 1, 2 and 3

$-(b_{13}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{38}^{''})^{(7,7,)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3

$-(b_{36}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)$ are eighth detrition coefficients for category 1, 2 and 3

$-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3

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$$\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \begin{bmatrix} (a_{44}^{'})^{(9)}\boxed{+(a_{44}^{''})^{(9)}(T_{45},t)} & \boxed{+(a_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)} \end{bmatrix} G_{13}$$

$$\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \begin{bmatrix} (a_{45}^{'})^{(9)}\boxed{+(a_{45}^{''})^{(9)}(T_{45},t)} & \boxed{+(a_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)} \end{bmatrix} G_{14}$$

$$\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \begin{bmatrix} (a_{46}^{'})^{(9)}\boxed{+(a_{46}^{''})^{(9)}(T_{37},t)} & \boxed{+(a_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)} \end{bmatrix} G_{15}$$

Where $+(a_{44}^{''})^{(9)}(T_{45},t)$, $+(a_{45}^{''})^{(9)}(T_{45},t)$, $+(a_{46}^{''})^{(9)}(T_{37},t)$ are first augmentation coefficients for category 1, 2 and 3

$+(a_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)$ are second augmentation coefficient for category 1, 2 and 3

$+(a_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3

$+(a_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3

$+(a_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficient for category 1, 2 and 3

$+(a_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$ are sixth

augmentation coefficient for category 1, 2 and 3
$+(a_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)$ are Seventh augmentation coefficient for category 1, 2 and 3
$+(a_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)$ $+(a_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)$ $+(a_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)$ are eighth augmentation coefficient for 1,2,3
$+(a_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{44}}{dt} = (b_{44})^{(9)}T_{45} - \begin{bmatrix} (b_{44}^{'})^{(9)}\left[-(b_{44}^{''})^{(9)}(G_{47},t)\right] & \left[-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)\right] & \left[-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)\right] \\ \left[-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)\right] & \left[-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)\right] & \left[-(b_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)\right] \\ \left[-(b_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)\right] & \left[-(b_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)\right] & \left[-(b_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)\right] \end{bmatrix} T_{13}$$

$$\frac{dT_{45}}{dt} = (b_{45})^{(9)}T_{44} - \begin{bmatrix} (b_{45}^{'})^{(9)}\left[-(b_{45}^{''})^{(9)}(G_{47},t)\right] & \left[-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)\right] & \left[-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)\right] \\ \left[-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)\right] & \left[-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)\right] & \left[-(b_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)\right] \\ \left[-(b_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)\right] & \left[-(b_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)\right] & \left[-(b_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)\right] \end{bmatrix} T_{14}$$

$$\frac{dT_{46}}{dt} = (b_{46})^{(9)}T_{45} - \begin{bmatrix} (b_{46}^{'})^{(9)}\left[-(b_{46}^{''})^{(9)}(G_{47},t)\right] & \left[-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)\right] & \left[-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)\right] \\ \left[-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)\right] & \left[-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)\right] & \left[-(b_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)\right] \\ \left[-(b_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)\right] & \left[-(b_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)\right] & \left[-(b_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)\right] \end{bmatrix} T_{15}$$

Where $-(b_{44}^{''})^{(9)}(G_{47},t)$, $-(b_{45}^{''})^{(9)}(G_{47},t)$, $-(b_{46}^{''})^{(9)}(G_{47},t)$ are first detrition coefficients for category 1, 2 and 3
$-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3
$-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3
$-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3
$-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3
$-(b_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)$ are sixth detrition coefficients for category 1, 2 and 3
$-(b_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)$ are seventh detrition coefficients for category 1, 2 and 3

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| $\boxed{-(b_{37}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{36}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{38}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$ are eighth detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$ are ninth detrition coefficients for category 1, 2 and 3 | |
| Where we suppose | |
| $(a_i)^{(1)}, (a_i^{'})^{(1)}, (a_i^{''})^{(1)}, (b_i)^{(1)}, (b_i^{'})^{(1)}, (b_i^{''})^{(1)} > 0, \quad i,j = 13,14,15$<br><br>The functions $(a_i^{''})^{(1)}, (b_i^{''})^{(1)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(1)}$, $(r_i)^{(1)}$:<br><br>$(a_i^{''})^{(1)}(T_{14},t) \leq (p_i)^{(1)} \leq (\hat{A}_{13})^{(1)}$<br><br>$(b_i^{''})^{(1)}(G,t) \leq (r_i)^{(1)} \leq (b_i^{'})^{(1)} \leq (\hat{B}_{13})^{(1)}$ | *97* |
| $lim_{T_2 \to \infty}(a_i^{''})^{(1)}(T_{14},t) = (p_i)^{(1)}$<br><br>$\lim_{G \to \infty}(b_i^{''})^{(1)}(G,t) = (r_i)^{(1)}$<br><br>**Definition of** $(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}$ :<br><br>Where $\boxed{(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}}$ are positive constants and $\boxed{i = 13,14,15}$ | 98 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(1)}(T_{14}^{'},t) - (a_i^{''})^{(1)}(T_{14},t)\rvert \leq (\hat{k}_{13})^{(1)}\lvert T_{14} - T_{14}^{'}\rvert e^{-(\hat{M}_{13})^{(1)}t}$<br><br>$\lvert(b_i^{''})^{(1)}(G^{'},t) - (b_i^{''})^{(1)}(G,t)\rvert < (\hat{k}_{13})^{(1)}\lvert\lvert G - G^{'}\rvert\rvert e^{-(\hat{M}_{13})^{(1)}t}$ | 99 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(1)}(T_{14}^{'},t)$ and $(a_i^{''})^{(1)}(T_{14},t)$ . $(T_{14}^{'},t)$ and $(T_{14},t)$ are points belonging to the interval $\left[(\hat{k}_{13})^{(1)}, (\hat{M}_{13})^{(1)}\right]$. It is to be noted that $(a_i^{''})^{(1)}(T_{14},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{13})^{(1)} = 1$ then the function $(a_i^{''})^{(1)}(T_{14},t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}$ :<br><br>$(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}$, are positive constants<br><br>$\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}} < 1$ | 100 |
| **Definition of** $(\hat{P}_{13})^{(1)}, (\hat{Q}_{13})^{(1)}$ :<br><br>There exists two constants $(\hat{P}_{13})^{(1)}$ and $(\hat{Q}_{13})^{(1)}$ which together With $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}, (\hat{A}_{13})^{(1)}$ and $(\hat{B}_{13})^{(1)}$ and the constants $(a_i)^{(1)}, (a_i^{'})^{(1)}, (b_i)^{(1)}, (b_i^{'})^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}, i = 13,14,15,$ | 101 |

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| satisfy the inequalities<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(a_i)^{(1)} + (a_i')^{(1)} + (\hat{A}_{13})^{(1)} + (\hat{P}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(b_i)^{(1)} + (b_i')^{(1)} + (\hat{B}_{13})^{(1)} + (\hat{Q}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$ | |
| Where we suppose | |
| $(a_i)^{(2)}, (a_i')^{(2)}, (a_i'')^{(2)}, (b_i)^{(2)}, (b_i')^{(2)}, (b_i'')^{(2)} > 0, \quad i,j = 16,17,18$ | |
| The functions $(a_i'')^{(2)}, (b_i'')^{(2)}$ are positive continuous increasing and bounded. | |
| **Definition of** $(p_i)^{(2)}$, $(r_i)^{(2)}$: | |
| $(a_i'')^{(2)}(T_{17}, t) \leq (p_i)^{(2)} \leq (\hat{A}_{16})^{(2)}$ | 102 |
| $(b_i'')^{(2)}(G_{19}, t) \leq (r_i)^{(2)} \leq (b_i')^{(2)} \leq (\hat{B}_{16})^{(2)}$ | 103 |
| $\lim_{T_2 \to \infty} (a_i'')^{(2)}(T_{17}, t) = (p_i)^{(2)}$ | 104 |
| $\lim_{G \to \infty} (b_i'')^{(2)}((G_{19}), t) = (r_i)^{(2)}$ | 105 |
| **Definition of** $(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}$ :<br>Where $(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}, (p_i)^{(2)}, (r_i)^{(2)}$ are positive constants and $i = 16,17,18$ | 106 |
| They satisfy Lipschitz condition: | |
| $\lvert(a_i'')^{(2)}(T_{17}', t) - (a_i'')^{(2)}(T_{17}, t)\rvert \leq (\hat{k}_{16})^{(2)}\lvert T_{17} - T_{17}'\rvert e^{-(\hat{M}_{16})^{(2)}t}$ | 107 |
| $\lvert(b_i'')^{(2)}((G_{19})', t) - (b_i'')^{(2)}((G_{19}), t)\rvert < (\hat{k}_{16})^{(2)}\lvert\lvert(G_{19}) - (G_{19})'\rvert\rvert e^{-(\hat{M}_{16})^{(2)}t}$ | 108 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(2)}(T_{17}', t)$ and $(a_i'')^{(2)}(T_{17}, t)$ . $(T_{17}', t)$ and $(T_{17}, t)$ are points belonging to the interval $[(\hat{k}_{16})^{(2)}, (\hat{M}_{16})^{(2)}]$ . It is to be noted that $(a_i'')^{(2)}(T_{17}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{16})^{(2)} = 1$ then the function $(a_i'')^{(2)}(T_{17}, t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$ : | |
| $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$, are positive constants<br>$\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}, \frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ | 109 |
| **Definition of** $(\hat{P}_{13})^{(2)}, (\hat{Q}_{13})^{(2)}$ :<br>There exists two constants $(\hat{P}_{16})^{(2)}$ and $(\hat{Q}_{16})^{(2)}$ which together | |

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| with $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}, (\hat{A}_{16})^{(2)} and\ (\hat{B}_{16})^{(2)}$ and the constants $(a_i)^{(2)}, (a_i^{'})^{(2)}, (b_i)^{(2)}, (b_i^{'})^{(2)}, (p_i)^{(2)}, (r_i)^{(2)}, i = 16,17,18,$ satisfy the inequalities | |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(a_i)^{(2)} + (a_i^{'})^{(2)} + (\hat{A}_{16})^{(2)} + (\hat{P}_{16})^{(2)} (\hat{k}_{16})^{(2)}] < 1$ | 110 |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(b_i)^{(2)} + (b_i^{'})^{(2)} + (\hat{B}_{16})^{(2)} + (\hat{Q}_{16})^{(2)} (\hat{k}_{16})^{(2)}] < 1$ | 111 |
| Where we suppose | |
| $(a_i)^{(3)}, (a_i^{'})^{(3)}, (a_i^{''})^{(3)}, (b_i)^{(3)}, (b_i^{'})^{(3)}, (b_i^{''})^{(3)} > 0, \quad i, j = 20,21,22$<br>The functions $(a_i^{''})^{(3)}, (b_i^{''})^{(3)}$ are positive continuous increasing and bounded.<br>**Definition of** $(p_i)^{(3)}$, $(r_i)^{(3)}$:<br>$(a_i^{''})^{(3)}(T_{21}, t) \leq (p_i)^{(3)} \leq (\hat{A}_{20})^{(3)}$<br>$(b_i^{''})^{(3)}(G_{23}, t) \leq (r_i)^{(3)} \leq (b_i^{'})^{(3)} \leq (\hat{B}_{20})^{(3)}$ | 112 |
| $lim_{T_2 \to \infty}(a_i^{''})^{(3)}(T_{21}, t) = (p_i)^{(3)}$<br>$\lim_{G \to \infty}(b_i^{''})^{(3)}(G_{23}, t) = (r_i)^{(3)}$<br>**Definition of** $(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}$ :<br>Where $\boxed{(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}}$ are positive constants and $\boxed{i = 20,21,22}$ | 113 |
| They satisfy Lipschitz condition:<br>$\lvert(a_i^{''})^{(3)}(T_{21}^{'}, t) - (a_i^{''})^{(3)}(T_{21}, t)\rvert \leq (\hat{k}_{20})^{(3)} \lvert T_{21} - T_{21}^{'}\rvert e^{-(\hat{M}_{20})^{(3)}t}$<br>$\lvert(b_i^{''})^{(3)}(G_{23}^{'}, t) - (b_i^{''})^{(3)}(G_{23}, t)\rvert < (\hat{k}_{20})^{(3)} \lvert\lvert G_{23} - G_{23}^{'}\rvert\rvert e^{-(\hat{M}_{20})^{(3)}t}$ | 114 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(3)}(T_{21}^{'}, t)$ and $(a_i^{''})^{(3)}(T_{21}, t)$ . $(T_{21}^{'}, t)$ And $(T_{21}, t)$ are points belonging to the interval $[(\hat{k}_{20})^{(3)}, (\hat{M}_{20})^{(3)}]$. It is to be noted that $(a_i^{''})^{(3)}(T_{21}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{20})^{(3)} = 1$ then the function $(a_i^{''})^{(3)}(T_{21}, t)$, the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$ :<br>$(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$, are positive constants<br>$\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}, \frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}} < 1$ | 115 |

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| There exists two constants There exists two constants $(\hat{P}_{20})^{(3)}$ and $(\hat{Q}_{20})^{(3)}$ which together with $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}, (\hat{A}_{20})^{(3)} and\ (\hat{B}_{20})^{(3)}$ and the constants $(a_i)^{(3)}, (a_i^{'})^{(3)}, (b_i)^{(3)}, (b_i^{'})^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}, i = 20,21,22,$ satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(a_i)^{(3)} + (a_i^{'})^{(3)} + (\hat{A}_{20})^{(3)} + (\hat{P}_{20})^{(3)} (\hat{k}_{20})^{(3)}] < 1$<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(b_i)^{(3)} + (b_i^{'})^{(3)} + (\hat{B}_{20})^{(3)} + (\hat{Q}_{20})^{(3)} (\hat{k}_{20})^{(3)}] < 1$ | 116 |
| Where we suppose | |
| $(a_i)^{(4)}, (a_i^{'})^{(4)}, (a_i^{''})^{(4)}, (b_i)^{(4)}, (b_i^{'})^{(4)}, (b_i^{''})^{(4)} > 0, \quad i,j = 24,25,26$<br><br>The functions $(a_i^{''})^{(4)}, (b_i^{''})^{(4)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(4)}$, $(r_i)^{(4)}$:<br><br>$(a_i^{''})^{(4)}(T_{25}, t) \leq (p_i)^{(4)} \leq (\hat{A}_{24})^{(4)}$<br><br>$(b_i^{''})^{(4)}((G_{27}), t) \leq (r_i)^{(4)} \leq (b_i^{'})^{(4)} \leq (\hat{B}_{24})^{(4)}$ | 117 |
| $lim_{T_2 \to \infty}(a_i^{''})^{(4)}(T_{25}, t) = (p_i)^{(4)}$<br><br>$\lim_{G \to \infty}(b_i^{''})^{(4)}((G_{27}), t) = (r_i)^{(4)}$<br><br>**Definition of** $(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}$ :<br><br>Where $\boxed{(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}, (p_i)^{(4)}, (r_i)^{(4)}}$ are positive constants and $\boxed{i = 24,25,26}$ | 118 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(4)}(T_{25}^{'}, t) - (a_i^{''})^{(4)}(T_{25}, t)\rvert \leq (\hat{k}_{24})^{(4)} \lvert T_{25} - T_{25}^{'}\rvert e^{-(\hat{M}_{24})^{(4)}t}$<br><br>$\lvert(b_i^{''})^{(4)}((G_{27})^{'}, t) - (b_i^{''})^{(4)}((G_{27}), t)\rvert < (\hat{k}_{24})^{(4)} \lvert\lvert(G_{27}) - (G_{27})^{'}\rvert\rvert e^{-(\hat{M}_{24})^{(4)}t}$ | 119 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(4)}(T_{25}^{'}, t)$ and $(a_i^{''})^{(4)}(T_{25}, t)$ . $(T_{25}^{'}, t)$ and $(T_{25}, t)$ are points belonging to the interval $[(\hat{k}_{24})^{(4)}, (\hat{M}_{24})^{(4)}]$. It is to be noted that $(a_i^{''})^{(4)}(T_{25}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{24})^{(4)} = 1$ then the function $(a_i^{''})^{(4)}(T_{25}, t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$ :<br><br>$(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$, are positive constants | 120 |

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| $\frac{(a_i)^{(4)}}{(\hat{M}_{24})^{(4)}}, \frac{(b_i)^{(4)}}{(\hat{M}_{24})^{(4)}} < 1$ | |
| **Definition of** $(\hat{P}_{24})^{(4)}, (\hat{Q}_{24})^{(4)}$ :<br><br>There exists two constants $(\hat{P}_{24})^{(4)}$ and $(\hat{Q}_{24})^{(4)}$ which together with $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}, (\hat{A}_{24})^{(4)} and\ (\hat{B}_{24})^{(4)}$ and the constants $(a_i)^{(4)}, (a_i^{'})^{(4)}, (b_i)^{(4)}, (b_i^{'})^{(4)}, (p_i)^{(4)}, (r_i)^{(4)}, i = 24,25,26$, satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{24})^{(4)}}[(a_i)^{(4)} + (a_i^{'})^{(4)} + (\hat{A}_{24})^{(4)} + (\hat{P}_{24})^{(4)} (\hat{k}_{24})^{(4)}] < 1$<br><br>$\frac{1}{(\hat{M}_{24})^{(4)}}[(b_i)^{(4)} + (b_i^{'})^{(4)} + (\hat{B}_{24})^{(4)} + (\hat{Q}_{24})^{(4)} (\hat{k}_{24})^{(4)}] < 1$ | 121 |
| Where we suppose | |
| $(a_i)^{(5)}, (a_i^{'})^{(5)}, (a_i^{''})^{(5)}, (b_i)^{(5)}, (b_i^{'})^{(5)}, (b_i^{''})^{(5)} > 0, \quad i,j = 28,29,30$<br><br>The functions $(a_i^{''})^{(5)}, (b_i^{''})^{(5)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(5)}, (r_i)^{(5)}$:<br><br>$(a_i^{''})^{(5)}(T_{29}, t) \leq (p_i)^{(5)} \leq (\hat{A}_{28})^{(5)}$<br><br>$(b_i^{''})^{(5)}((G_{31}), t) \leq (r_i)^{(5)} \leq (b_i^{'})^{(5)} \leq (\hat{B}_{28})^{(5)}$ | 122 |
| $lim_{T_2 \to \infty} (a_i^{''})^{(5)} (T_{29}, t) = (p_i)^{(5)}$<br><br>$\lim_{G \to \infty} (b_i^{''})^{(5)} (G_{31}, t) = (r_i)^{(5)}$<br><br>**Definition of** $(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}$ :<br><br>Where $\boxed{(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}, (p_i)^{(5)}, (r_i)^{(5)}}$ are positive constants and $\boxed{i = 28,29,30}$ | 123 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(5)}(T_{29}^{'}, t) - (a_i^{''})^{(5)}(T_{29}, t)\rvert \leq (\hat{k}_{28})^{(5)} \lvert T_{29} - T_{29}^{'}\rvert e^{-(\hat{M}_{28})^{(5)} t}$<br><br>$\lvert(b_i^{''})^{(5)}((G_{31})', t) - (b_i^{''})^{(5)}((G_{31}), t)\rvert < (\hat{k}_{28})^{(5)} \lvert\lvert(G_{31}) - (G_{31})'\rvert\rvert e^{-(\hat{M}_{28})^{(5)} t}$ | 124 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(5)}(T_{29}^{'}, t)$ and $(a_i^{''})^{(5)}(T_{29}, t)$ . $(T_{29}^{'}, t)$ and $(T_{29}, t)$ are points belonging to the interval $[(\hat{k}_{28})^{(5)}, (\hat{M}_{28})^{(5)}]$. It is to be noted that $(a_i^{''})^{(5)}(T_{29}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{28})^{(5)} = 1$ then the function $(a_i^{''})^{(5)}(T_{29}, t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}$ : | 125 |

$(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}$, are positive constants

$$\frac{(a_i)^{(5)}}{(\hat{M}_{28})^{(5)}}, \frac{(b_i)^{(5)}}{(\hat{M}_{28})^{(5)}} < 1$$

**Definition of** $(\hat{P}_{28})^{(5)}, (\hat{Q}_{28})^{(5)}$ : (126)

There exists two constants $(\hat{P}_{28})^{(5)}$ and $(\hat{Q}_{28})^{(5)}$ which together with $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}, (\hat{A}_{28})^{(5)} and\ (\hat{B}_{28})^{(5)}$ and the constants $(a_i)^{(5)}, (a_i')^{(5)}, (b_i)^{(5)}, (b_i')^{(5)}, (p_i)^{(5)}, (r_i)^{(5)}, i = 28,29,30,$ satisfy the inequalities

$$\frac{1}{(\hat{M}_{28})^{(5)}}[(a_i)^{(5)} + (a_i')^{(5)} + (\hat{A}_{28})^{(5)} + (\hat{P}_{28})^{(5)} (\hat{k}_{28})^{(5)}] < 1$$

$$\frac{1}{(\hat{M}_{28})^{(5)}}[(b_i)^{(5)} + (b_i')^{(5)} + (\hat{B}_{28})^{(5)} + (\hat{Q}_{28})^{(5)} (\hat{k}_{28})^{(5)}] < 1$$

Where we suppose

$(a_i)^{(6)}, (a_i')^{(6)}, (a_i'')^{(6)}, (b_i)^{(6)}, (b_i')^{(6)}, (b_i'')^{(6)} > 0, \quad i, j = 32,33,34$ (127)

The functions $(a_i'')^{(6)}, (b_i'')^{(6)}$ are positive continuous increasing and bounded.

**Definition of** $(p_i)^{(6)}$, $(r_i)^{(6)}$:

$$(a_i'')^{(6)}(T_{33}, t) \le (p_i)^{(6)} \le (\hat{A}_{32})^{(6)}$$

$$(b_i'')^{(6)}((G_{35}), t) \le (r_i)^{(6)} \le (b_i')^{(6)} \le (\hat{B}_{32})^{(6)}$$

$$lim_{T_2 \to \infty} (a_i'')^{(6)} (T_{33}, t) = (p_i)^{(6)}$$ (128)

$$\lim_{G \to \infty} (b_i'')^{(6)} ((G_{35}), t) = (r_i)^{(6)}$$

**Definition of** $(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}$ :

Where $\boxed{(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}, (p_i)^{(6)}, (r_i)^{(6)}}$ are positive constants and $\boxed{i = 32,33,34}$

They satisfy Lipschitz condition:

$$|(a_i'')^{(6)}(T_{33}', t) - (a_i'')^{(6)}(T_{33}, t)| \le (\hat{k}_{32})^{(6)} |T_{33} - T_{33}'| e^{-(\hat{M}_{32})^{(6)} t}$$

$$|(b_i'')^{(6)}((G_{35})', t) - (b_i'')^{(6)}((G_{35}), t)| < (\hat{k}_{32})^{(6)} ||(G_{35}) - (G_{35})'|| e^{-(\hat{M}_{32})^{(6)} t}$$

With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(6)}(T_{33}', t)$ and $(a_i'')^{(6)}(T_{33}, t)$ . $(T_{33}', t)$ and $(T_{33}, t)$ are points belonging to the interval $[(\hat{k}_{32})^{(6)}, (\hat{M}_{32})^{(6)}]$. It is to be noted that $(a_i'')^{(6)}(T_{33}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{32})^{(6)} = 1$ then the function $(a_i'')^{(6)}(T_{33}, t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

| | |
|---|---|
| **Definition of** $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}$ :<br><br>$(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)},$ are positive constants<br><br>$\frac{(a_i)^{(6)}}{(\hat{M}_{32})^{(6)}}$ , $\frac{(b_i)^{(6)}}{(\hat{M}_{32})^{(6)}} < 1$ | 129 |
| **Definition of** $(\hat{P}_{32})^{(6)}, (\hat{Q}_{32})^{(6)}$ :<br><br>There exists two constants $(\hat{P}_{32})^{(6)}$ and $(\hat{Q}_{32})^{(6)}$ which together with $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}, (\hat{A}_{32})^{(6)} and\ (\hat{B}_{32})^{(6)}$ and the constants $(a_i)^{(6)}, (a_i^{'})^{(6)}, (b_i)^{(6)}, (b_i^{'})^{(6)}, (p_i)^{(6)}, (r_i)^{(6)}, i = 32,33,34,$ satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{32})^{(6)}}[(a_i)^{(6)} + (a_i^{'})^{(6)} + (\hat{A}_{32})^{(6)} + (\hat{P}_{32})^{(6)} (\hat{k}_{32})^{(6)}] < 1$<br><br>$\frac{1}{(\hat{M}_{32})^{(6)}}[(b_i)^{(6)} + (b_i^{'})^{(6)} + (\hat{B}_{32})^{(6)} + (\hat{Q}_{32})^{(6)} (\hat{k}_{32})^{(6)}] < 1$ | 130 |
| Where we suppose | |
| (Y) $(a_i)^{(7)}, (a_i^{'})^{(7)}, (a_i^{''})^{(7)}, (b_i)^{(7)}, (b_i^{'})^{(7)}, (b_i^{''})^{(7)} > 0, \quad i,j = 36,37,38$<br><br>(Z) The functions $(a_i^{''})^{(7)}, (b_i^{''})^{(7)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(7)}$, $(r_i)^{(7)}$:<br><br>$(a_i^{''})^{(7)}(T_{37}, t) \leq (p_i)^{(7)} \leq (\hat{A}_{36})^{(7)}$<br><br>$(b_i^{''})^{(7)}(G_{39}, t) \leq (r_i)^{(7)} \leq (b_i^{'})^{(7)} \leq (\hat{B}_{36})^{(7)}$ | 131 |
| *(AA)* $\lim_{T_2 \to \infty} (a_i^{''})^{(7)}(T_{37}, t) = (p_i)^{(7)}$<br>*(BB)*<br>$\lim_{G \to \infty} (b_i^{''})^{(7)}((G_{39}), t) = (r_i)^{(7)}$<br><br>**Definition of** $(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}$ :<br><br>Where $\boxed{(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}, (p_i)^{(7)}, (r_i)^{(7)}}$ are positive constants and $\boxed{i = 36,37,38}$ | 132 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(7)}(T_{37}^{'}, t) - (a_i^{''})^{(7)}(T_{37}, t)\rvert \leq (\hat{k}_{36})^{(7)} \lvert T_{37} - T_{37}^{'}\rvert e^{-(\hat{M}_{36})^{(7)}t}$<br><br>$\lvert(b_i^{''})^{(7)}((G_{39})^{'}, t) - (b_i^{''})^{(7)}((G_{39}), t)\rvert < (\hat{k}_{36})^{(7)} \lvert\lvert(G_{39}) - (G_{39})^{'}\rvert\rvert e^{-(\hat{M}_{36})^{(7)}t}$ | 133 |

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|---|---|
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(7)}(T_{37}^{'},t)$ and$(a_i^{''})^{(7)}(T_{37},t)$ . $(T_{37}^{'},t)$ and $(T_{37},t)$ are points belonging to the interval $\left[(\hat{k}_{36})^{(7)},(\widehat{M}_{36})^{(7)}\right]$. It is to be noted that $(a_i^{''})^{(7)}(T_{37},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\widehat{M}_{36})^{(7)}=1$ then the function $(a_i^{''})^{(7)}(T_{37},t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\widehat{M}_{36})^{(7)},(\hat{k}_{36})^{(7)}$ : <br><br>(CC) $(\widehat{M}_{36})^{(7)},(\hat{k}_{36})^{(7)},$ are positive constants <br><br>$\frac{(a_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}$ , $\frac{(b_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}<1$ | 134 |
| **Definition of** $(\hat{P}_{36})^{(7)},(\hat{Q}_{36})^{(7)}$ : <br><br>(DD) There exists two constants $(\hat{P}_{36})^{(7)}$ and $(\hat{Q}_{36})^{(7)}$ which together with $(\widehat{M}_{36})^{(7)},(\hat{k}_{36})^{(7)},(\hat{A}_{36})^{(7)} and\ (\hat{B}_{36})^{(7)}$ and the constants $(a_i)^{(7)},(a_i^{'})^{(7)},(b_i)^{(7)},(b_i^{'})^{(7)},(p_i)^{(7)},\ (r_i)^{(7)},i=36,37,38$, satisfy the inequalities <br><br>$\frac{1}{(\widehat{M}_{36})^{(7)}}[(a_i)^{(7)}+(a_i^{'})^{(7)}+(\hat{A}_{36})^{(7)}+(\hat{P}_{36})^{(7)}(\hat{k}_{36})^{(7)}]<1$ <br><br>$\frac{1}{(\widehat{M}_{36})^{(7)}}[(b_i)^{(7)}+(b_i^{'})^{(7)}+(\hat{B}_{36})^{(7)}+(\hat{Q}_{36})^{(7)}(\hat{k}_{36})^{(7)}]<1$ | 135 |
| Where we suppose | |
| $(a_i)^{(8)},(a_i^{'})^{(8)},(a_i^{''})^{(8)},(b_i)^{(8)},(b_i^{'})^{(8)},(b_i^{''})^{(8)}>0,\quad i,j=40,41,42$ | 136 |
| The functions $(a_i^{''})^{(8)},(b_i^{''})^{(8)}$ are positive continuous increasing and bounded | |
| **Definition of** $(p_i)^{(8)},\ (r_i)^{(8)}$: | 137 |
| $(a_i^{''})^{(8)}(T_{41},t)\leq(p_i)^{(8)}\leq(\hat{A}_{40})^{(8)}$ | 138 |
| $(b_i^{''})^{(8)}((G_{43}),t)\leq\ (r_i)^{(8)}\leq(b_i^{'})^{(8)}\leq(\hat{B}_{40})^{(8)}$ | 139 |
| $\lim_{T_2\to\infty}(a_i^{''})^{(8)}(T_{41},t)=(p_i)^{(8)}$ | 140 |
| $\lim_{G\to\infty}(b_i^{''})^{(8)}\left((G_{43}),t\right)=\ (r_i)^{(8)}$ | 141 |
| **Definition of** $(\hat{A}_{40})^{(8)},(\hat{B}_{40})^{(8)}$ : <br><br>Where $\boxed{(\hat{A}_{40})^{(8)},(\hat{B}_{40})^{(8)},(p_i)^{(8)},\ (r_i)^{(8)}}$ are positive constants and $\boxed{i=40,41,42}$ | |

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| They satisfy Lipschitz condition: | |
| $\|(a_i^{''})^{(8)}(T_{41}^{'},t)-(a_i^{''})^{(8)}(T_{41},t)\| \leq (\hat{k}_{40})^{(8)}\|T_{41}-T_{41}^{'}\|e^{-(\hat{M}_{40})^{(8)}t}$ | 142 |
| $\|(b_i^{''})^{(8)}((G_{43})^{'},t)-(b_i^{''})^{(8)}((G_{43}),t)\| < (\hat{k}_{40})^{(8)}\|\|(G_{43})-(G_{43})^{'}\|\|e^{-(\hat{M}_{40})^{(8)}t}$ | 143 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(8)}(T_{41}^{'},t)$ and $(a_i^{''})^{(8)}(T_{41},t)$ . $(T_{41}^{'},t)$ and $(T_{41},t)$ are points belonging to the interval $[(\hat{k}_{40})^{(8)},(\hat{M}_{40})^{(8)}]$. It is to be noted that $(a_i^{''})^{(8)}(T_{41},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{40})^{(8)}=1$ then the function $(a_i^{''})^{(8)}(T_{41},t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{40})^{(8)},(\hat{k}_{40})^{(8)}$ : | |
| $(\hat{M}_{40})^{(8)},(\hat{k}_{40})^{(8)}$, are positive constants | |
| $\frac{(a_i)^{(8)}}{(\hat{M}_{40})^{(8)}}$ , $\frac{(b_i)^{(8)}}{(\hat{M}_{40})^{(8)}}<1$ | 144 |
| **Definition of** $(\hat{P}_{40})^{(8)},(\hat{Q}_{40})^{(8)}$ :<br><br>There exists two constants $(\hat{P}_{40})^{(8)}$ and $(\hat{Q}_{40})^{(8)}$ which together with $(\hat{M}_{40})^{(8)},(\hat{k}_{40})^{(8)},(\hat{A}_{40})^{(8)}$ $(\hat{B}_{40})^{(8)}$ and the constants $(a_i)^{(8)},(a_i^{'})^{(8)},(b_i)^{(8)},(b_i^{'})^{(8)},(p_i)^{(8)}$, $(r_i)^{(8)}, i=40,41,42,$<br>Satisfy the inequalities | |
| $\frac{1}{(\hat{M}_{40})^{(8)}}[(a_i)^{(8)}+(a_i^{'})^{(8)}+(\hat{A}_{40})^{(8)}+(\hat{P}_{40})^{(8)}(\hat{k}_{40})^{(8)}]<1$ | 145 |
| $\frac{1}{(\hat{M}_{40})^{(8)}}[(b_i)^{(8)}+(b_i^{'})^{(8)}+(\hat{B}_{40})^{(8)}+(\hat{Q}_{40})^{(8)}(\hat{k}_{40})^{(8)}]<1$ | 146 |
| Where we suppose | |
| $(a_i)^{(9)},(a_i^{'})^{(9)},(a_i^{''})^{(9)},(b_i)^{(9)},(b_i^{'})^{(9)},(b_i^{''})^{(9)}>0, \quad i,j=44,45,46$<br><br>The functions $(a_i^{''})^{(9)},(b_i^{''})^{(9)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(9)}$, $(r_i)^{(9)}$:<br><br>$(a_i^{''})^{(9)}(T_{45},t) \leq (p_i)^{(9)} \leq (\hat{A}_{44})^{(9)}$<br><br>$(b_i^{''})^{(9)}(G_{47},t) \leq (r_i)^{(9)} \leq (b_i^{'})^{(9)} \leq (\hat{B}_{44})^{(9)}$ | 146 A |
| $\lim_{T_2\to\infty}(a_i^{''})^{(9)}(T_{45},t)=(p_i)^{(9)}$<br><br>$\lim_{G\to\infty}(b_i^{''})^{(9)}(G_{47},t)=(r_i)^{(9)}$ | |

**Definition of** $(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}$ :

Where $\boxed{(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}, (p_i)^{(9)}, \ (r_i)^{(9)}}$ are positive constants and $\boxed{i = 44{,}45{,}46}$

They satisfy Lipschitz condition:

$$|(a_i^{''})^{(9)}(T_{45}^{'}, t) - (a_i^{''})^{(9)}(T_{45}, t)| \leq (\hat{k}_{44})^{(9)}|T_{45} - T_{45}^{'}|e^{-(\hat{M}_{44})^{(9)}t}$$

$$|(b_i^{''})^{(9)}((G_{47})^{'}, t) - (b_i^{''})^{(9)}\big((G_{47}), t\big)| < (\hat{k}_{44})^{(9)}||(G_{47}) - (G_{47})^{'}||e^{-(\hat{M}_{44})^{(9)}t}$$

With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(9)}(T_{45}^{'}, t)$ and$(a_i^{''})^{(9)}(T_{45}, t)$ . $(T_{45}^{'}, t)$ and $(T_{45}, t)$ are points belonging to the interval $\left[(\hat{k}_{44})^{(9)}, (\hat{M}_{44})^{(9)}\right]$. It is to be noted that $(a_i^{''})^{(9)}(T_{45}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{44})^{(9)} = 1$ then the function $(a_i^{''})^{(9)}(T_{45}, t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

**Definition of** $(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}$ :

$(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)},$ are positive constants

$$\frac{(a_i)^{(9)}}{(\hat{M}_{44})^{(9)}} \ , \frac{(b_i)^{(9)}}{(\hat{M}_{44})^{(9)}} < 1$$

**Definition of** $(\hat{P}_{44})^{(9)}, (\hat{Q}_{44})^{(9)}$ :

There exists two constants $(\hat{P}_{44})^{(9)}$ and $(\hat{Q}_{44})^{(9)}$ which together with $(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}, (\hat{A}_{44})^{(9)} and \ (\hat{B}_{44})^{(9)}$ and the constants $(a_i)^{(9)}, (a_i^{'})^{(9)}, (b_i)^{(9)}, (b_i^{'})^{(9)}, (p_i)^{(9)}, \ (r_i)^{(9)}, i = 44{,}45{,}46,$ satisfy the inequalities

$$\frac{1}{(\hat{M}_{44})^{(9)}}[\,(a_i)^{(9)} + (a_i^{'})^{(9)} + \ (\hat{A}_{44})^{(9)} + (\hat{P}_{44})^{(9)}\,(\hat{k}_{44})^{(9)}] < 1$$

$$\frac{1}{(\hat{M}_{44})^{(9)}}[\ (b_i)^{(9)} + (b_i^{'})^{(9)} + \ (\hat{B}_{44})^{(9)} + (\hat{Q}_{44})^{(9)}\ (\hat{k}_{44})^{(9)}] < 1$$

**Theorem 1:** if the conditions above are fulfilled, there exists a solution satisfying the conditions 147

Definition of $G_i(0), T_i(0)$ :

$$G_i(t) \leq \left(\hat{P}_{13}\right)^{(1)} e^{(\hat{M}_{13})^{(1)}t} \ , \ \boxed{G_i(0) = G_i^0 > 0}$$

$$T_i(t) \leq (\hat{Q}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)}t} \ , \ \boxed{T_i(0) = T_i^0 > 0}$$

**Theorem 2 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions 148

**Definition of** $G_i(0), T_i(0)$

$$G_i(t) \leq (\hat{P}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)}t} \ , \ G_i(0) = G_i^0 > 0$$

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| $T_i(t) \leq (\hat{Q}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ , $T_i(0) = T_i^0 > 0$ | |
| **Theorem 3 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>$G_i(t) \leq (\hat{P}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ , $G_i(0) = G_i^0 > 0$<br>$T_i(t) \leq (\hat{Q}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ , $T_i(0) = T_i^0 > 0$ | 149 |
| **Theorem 4 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \leq (\hat{P}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \leq (\hat{Q}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 150 |
| **Theorem 5 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \leq (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \leq (\hat{Q}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 151 |
| **Theorem 6 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \leq (\hat{P}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \leq (\hat{Q}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 152 |
| **Theorem 7:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \leq (\hat{P}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \leq (\hat{Q}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 |
| **Theorem 8:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \leq (\hat{P}_{40})^{(8)} e^{(\hat{M}_{40})^{(8)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$ | 153 A |

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|---|---|
| $T_i(t) \leq (\hat{Q}_{40})^{(8)} e^{(\hat{M}_{40})^{(8)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | |
| **Theorem 9:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br><br>**Definition of** $G_i(0), T_i(0)$ :<br><br>$G_i(t) \leq (\hat{P}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br><br>$T_i(t) \leq (\hat{Q}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 B |
| **Proof:** Consider operator $\mathcal{A}^{(1)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 154 |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{13})^{(1)}$ , $T_i^0 \leq (\hat{Q}_{13})^{(1)}$, | 155 |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ | 156 |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ | 157 |
| By<br><br>$\bar{G}_{13}(t) = G_{13}^0 + \int_0^t \left[ (a_{13})^{(1)} G_{14}(s_{(13)}) - \left( (a_{13}')^{(1)} + a_{13}'')^{(1)} (T_{14}(s_{(13)}), s_{(13)}) \right) G_{13}(s_{(13)}) \right] ds_{(13)}$ | 158 |
| $\bar{G}_{14}(t) = G_{14}^0 + \int_0^t \left[ (a_{14})^{(1)} G_{13}(s_{(13)}) - \left( (a_{14}')^{(1)} + (a_{14}'')^{(1)} (T_{14}(s_{(13)}), s_{(13)}) \right) G_{14}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{G}_{15}(t) = G_{15}^0 + \int_0^t \left[ (a_{15})^{(1)} G_{14}(s_{(13)}) - \left( (a_{15}')^{(1)} + (a_{15}'')^{(1)} (T_{14}(s_{(13)}), s_{(13)}) \right) G_{15}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{T}_{13}(t) = T_{13}^0 + \int_0^t \left[ (b_{13})^{(1)} T_{14}(s_{(13)}) - \left( (b_{13}')^{(1)} - (b_{13}'')^{(1)} (G(s_{(13)}), s_{(13)}) \right) T_{13}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{T}_{14}(t) = T_{14}^0 + \int_0^t \left[ (b_{14})^{(1)} T_{13}(s_{(13)}) - \left( (b_{14}')^{(1)} - (b_{14}'')^{(1)} (G(s_{(13)}), s_{(13)}) \right) T_{14}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{T}_{15}(t) = T_{15}^0 + \int_0^t \left[ (b_{15})^{(1)} T_{14}(s_{(13)}) - \left( (b_{15}')^{(1)} - (b_{15}'')^{(1)} (G(s_{(13)}), s_{(13)}) \right) T_{15}(s_{(13)}) \right] ds_{(13)}$<br><br>Where $s_{(13)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br><br>Consider operator $\mathcal{A}^{(2)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 159 |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{16})^{(2)}$ , $T_i^0 \leq (\hat{Q}_{16})^{(2)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ | |

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| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{16})^{(2)} e^{(\bar{M}_{16})^{(2)} t}$ | |
| By<br>$\bar{G}_{16}(t) = G_{16}^0 + \int_0^t \left[ (a_{16})^{(2)} G_{17}(s_{(16)}) - \left( (a_{16}')^{(2)} + a_{16}'')^{(2)} (T_{17}(s_{(16)}), s_{(16)}) \right) G_{16}(s_{(16)}) \right] ds_{(16)}$ | 160 |
| $\bar{G}_{17}(t) = G_{17}^0 + \int_0^t \left[ (a_{17})^{(2)} G_{16}(s_{(16)}) - \left( (a_{17}')^{(2)} + (a_{17}'')^{(2)} (T_{17}(s_{(16)}), s_{(17)}) \right) G_{17}(s_{(16)}) \right] ds_{(16)}$ | |
| $\bar{G}_{18}(t) = G_{18}^0 + \int_0^t \left[ (a_{18})^{(2)} G_{17}(s_{(16)}) - \left( (a_{18}')^{(2)} + (a_{18}'')^{(2)} (T_{17}(s_{(16)}), s_{(16)}) \right) G_{18}(s_{(16)}) \right] ds_{(16)}$ | |
| $\bar{T}_{16}(t) = T_{16}^0 + \int_0^t \left[ (b_{16})^{(2)} T_{17}(s_{(16)}) - \left( (b_{16}')^{(2)} - (b_{16}'')^{(2)} (G_{19}(s_{(16)}), s_{(16)}) \right) T_{16}(s_{(16)}) \right] ds_{(16)}$ | |
| $\bar{T}_{17}(t) = T_{17}^0 + \int_0^t \left[ (b_{17})^{(2)} T_{16}(s_{(16)}) - \left( (b_{17}')^{(2)} - (b_{17}'')^{(2)} (G_{19}(s_{(16)}), s_{(16)}) \right) T_{17}(s_{(16)}) \right] ds_{(16)}$ | |
| $\bar{T}_{18}(t) = T_{18}^0 + \int_0^t \left[ (b_{18})^{(2)} T_{17}(s_{(16)}) - \left( (b_{18}')^{(2)} - (b_{18}'')^{(2)} (G_{19}(s_{(16)}), s_{(16)}) \right) T_{18}(s_{(16)}) \right] ds_{(16)}$<br>Where $s_{(16)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(3)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{20})^{(3)}$, $T_i^0 \le (\hat{Q}_{20})^{(3)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{20})^{(3)} e^{(\bar{M}_{20})^{(3)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{20})^{(3)} e^{(\bar{M}_{20})^{(3)} t}$ | |
| By<br>$\bar{G}_{20}(t) = G_{20}^0 + \int_0^t \left[ (a_{20})^{(3)} G_{21}(s_{(20)}) - \left( (a_{20}')^{(3)} + a_{20}'')^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{20}(s_{(20)}) \right] ds_{(20)}$ | 161 |
| $\bar{G}_{21}(t) = G_{21}^0 + \int_0^t \left[ (a_{21})^{(3)} G_{20}(s_{(20)}) - \left( (a_{21}')^{(3)} + (a_{21}'')^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{21}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{G}_{22}(t) = G_{22}^0 + \int_0^t \left[ (a_{22})^{(3)} G_{21}(s_{(20)}) - \left( (a_{22}')^{(3)} + (a_{22}'')^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{22}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{20}(t) = T_{20}^0 + \int_0^t \left[ (b_{20})^{(3)} T_{21}(s_{(20)}) - \left( (b_{20}')^{(3)} - (b_{20}'')^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{20}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{21}(t) = T_{21}^0 + \int_0^t \left[ (b_{21})^{(3)} T_{20}(s_{(20)}) - \left( (b_{21}')^{(3)} - (b_{21}'')^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{21}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{22}(t) = T_{22}^0 + \int_0^t \left[ (b_{22})^{(3)} T_{21}(s_{(20)}) - \left( (b_{22}')^{(3)} - (b_{22}'')^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{22}(s_{(20)}) \right] ds_{(20)}$<br>Where $s_{(20)}$ is the integrand that is integrated over an interval $(0, t)$ | |

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| **Proof:** Consider operator $\mathcal{A}^{(4)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i$: $\mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{24})^{(4)}$ , $T_i^0 \leq (\hat{Q}_{24})^{(4)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ | |
| By<br>$\bar{G}_{24}(t) = G_{24}^0 + \int_0^t \left[ (a_{24})^{(4)} G_{25}(s_{(24)}) - \left( (a'_{24})^{(4)} + a''_{24})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{24}(s_{(24)}) \right] ds_{(24)}$ | 162 |
| $\bar{G}_{25}(t) = G_{25}^0 + \int_0^t \left[ (a_{25})^{(4)} G_{24}(s_{(24)}) - \left( (a'_{25})^{(4)} + (a''_{25})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{25}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{G}_{26}(t) = G_{26}^0 + \int_0^t \left[ (a_{26})^{(4)} G_{25}(s_{(24)}) - \left( (a'_{26})^{(4)} + (a''_{26})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{26}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{24}(t) = T_{24}^0 + \int_0^t \left[ (b_{24})^{(4)} T_{25}(s_{(24)}) - \left( (b'_{24})^{(4)} - (b''_{24})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{24}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{25}(t) = T_{25}^0 + \int_0^t \left[ (b_{25})^{(4)} T_{24}(s_{(24)}) - \left( (b'_{25})^{(4)} - (b''_{25})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{25}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{26}(t) = T_{26}^0 + \int_0^t \left[ (b_{26})^{(4)} T_{25}(s_{(24)}) - \left( (b'_{26})^{(4)} - (b''_{26})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{26}(s_{(24)}) \right] ds_{(24)}$<br>Where $s_{(24)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:** Consider operator $\mathcal{A}^{(5)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i$: $\mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{28})^{(5)}$ , $T_i^0 \leq (\hat{Q}_{28})^{(5)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| By<br>$\bar{G}_{28}(t) = G_{28}^0 + \int_0^t \left[ (a_{28})^{(5)} G_{29}(s_{(28)}) - \left( (a'_{28})^{(5)} + a''_{28})^{(5)} (T_{29}(s_{(28)}), s_{(28)}) \right) G_{28}(s_{(28)}) \right] ds_{(28)}$ | 163 |
| $\bar{G}_{29}(t) = G_{29}^0 + \int_0^t \left[ (a_{29})^{(5)} G_{28}(s_{(28)}) - \left( (a'_{29})^{(5)} + (a''_{29})^{(5)} (T_{29}(s_{(28)}), s_{(28)}) \right) G_{29}(s_{(28)}) \right] ds_{(28)}$ | |
| $\bar{G}_{30}(t) = G_{30}^0 + \int_0^t \left[ (a_{30})^{(5)} G_{29}(s_{(28)}) - \left( (a'_{30})^{(5)} + (a''_{30})^{(5)} (T_{29}(s_{(28)}), s_{(28)}) \right) G_{30}(s_{(28)}) \right] ds_{(28)}$ | |
| $\bar{T}_{28}(t) = T_{28}^0 + \int_0^t \left[ (b_{28})^{(5)} T_{29}(s_{(28)}) - \left( (b'_{28})^{(5)} - (b''_{28})^{(5)} (G_{31}(s_{(28)}), s_{(28)}) \right) T_{28}(s_{(28)}) \right] ds_{(28)}$ | |

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| $\bar{T}_{29}(t) = T_{29}^0 + \int_0^t \left[(b_{29})^{(5)} T_{28}(s_{(28)}) - \left((b'_{29})^{(5)} - (b''_{29})^{(5)}(G_{31}(s_{(28)}), s_{(28)})\right) T_{29}(s_{(28)})\right] ds_{(28)}$ | |
| $\bar{\mathrm{T}}_{30}(\mathrm{t}) = \mathrm{T}_{30}^0 + \int_0^t \left[(b_{30})^{(5)} T_{29}(s_{(28)}) - \left((b'_{30})^{(5)} - (b''_{30})^{(5)}(G_{31}(s_{(28)}), s_{(28)})\right) T_{30}(s_{(28)})\right] ds_{(28)}$<br>Where $s_{(28)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(6)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \le (\hat{P}_{32})^{(6)}$ , $T_i^0 \le (\hat{Q}_{32})^{(6)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ | |
| By<br>$\bar{G}_{32}(t) = G_{32}^0 + \int_0^t \left[(a_{32})^{(6)} G_{33}(s_{(32)}) - \left((a'_{32})^{(6)} + a''_{32})^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{32}(s_{(32)})\right] ds_{(32)}$ | 164 |
| $\bar{G}_{33}(t) = G_{33}^0 + \int_0^t \left[(a_{33})^{(6)} G_{32}(s_{(32)}) - \left((a'_{33})^{(6)} + (a''_{33})^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{33}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{G}_{34}(t) = G_{34}^0 + \int_0^t \left[(a_{34})^{(6)} G_{33}(s_{(32)}) - \left((a'_{34})^{(6)} + (a''_{34})^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{34}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{T}_{32}(t) = T_{32}^0 + \int_0^t \left[(b_{32})^{(6)} T_{33}(s_{(32)}) - \left((b'_{32})^{(6)} - (b''_{32})^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{32}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{T}_{33}(t) = T_{33}^0 + \int_0^t \left[(b_{33})^{(6)} T_{32}(s_{(32)}) - \left((b'_{33})^{(6)} - (b''_{33})^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{33}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{\mathrm{T}}_{34}(\mathrm{t}) = \mathrm{T}_{34}^0 + \int_0^t \left[(b_{34})^{(6)} T_{33}(s_{(32)}) - \left((b'_{34})^{(6)} - (b''_{34})^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{34}(s_{(32)})\right] ds_{(32)}$<br>Where $s_{(32)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(7)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \le (\hat{P}_{36})^{(7)}$ , $T_i^0 \le (\hat{Q}_{36})^{(7)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ | |

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| By<br>$\bar{G}_{36}(t) = G_{36}^{0} + \int_0^t \left[(a_{36})^{(7)} G_{37}(s_{(36)}) - \left((a_{36}^{'})^{(7)} + a_{36}^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{36}(s_{(36)})\right] ds_{(36)}$ | 165 |
| $\bar{G}_{37}(t) = G_{37}^{0} + \int_0^t \left[(a_{37})^{(7)} G_{36}(s_{(36)}) - \left((a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{37}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{G}_{38}(t) = G_{38}^{0} + \int_0^t \left[(a_{38})^{(7)} G_{37}(s_{(36)}) - \left((a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{38}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{36}(t) = T_{36}^{0} + \int_0^t \left[(b_{36})^{(7)} T_{37}(s_{(36)}) - \left((b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{36}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{37}(t) = T_{37}^{0} + \int_0^t \left[(b_{37})^{(7)} T_{36}(s_{(36)}) - \left((b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{37}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{38}(t) = T_{38}^{0} + \int_0^t \left[(b_{38})^{(7)} T_{37}(s_{(36)}) - \left((b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{38}(s_{(36)})\right] ds_{(36)}$<br>Where $s_{(36)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(8)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{40})^{(8)}$, $T_i^0 \le (\hat{Q}_{40})^{(8)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{40})^{(8)} e^{(\bar{M}_{40})^{(8)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{40})^{(8)} e^{(\bar{M}_{40})^{(8)} t}$ | |
| By<br>$\bar{G}_{40}(t) = G_{40}^{0} + \int_0^t \left[(a_{40})^{(8)} G_{41}(s_{(40)}) - \left((a_{40}^{'})^{(8)} + a_{40}^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\right) G_{40}(s_{(40)})\right] ds_{(40)}$ | 166 |
| $\bar{G}_{41}(t) = G_{41}^{0} + \int_0^t \left[(a_{41})^{(8)} G_{40}(s_{(40)}) - \left((a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\right) G_{41}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{G}_{42}(t) = G_{42}^{0} + \int_0^t \left[(a_{42})^{(8)} G_{41}(s_{(40)}) - \left((a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\right) G_{42}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{T}_{40}(t) = T_{40}^{0} + \int_0^t \left[(b_{40})^{(8)} T_{41}(s_{(40)}) - \left((b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}(G_{43}(s_{(40)}), s_{(40)})\right) T_{40}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{T}_{41}(t) = T_{41}^{0} + \int_0^t \left[(b_{41})^{(8)} T_{40}(s_{(40)}) - \left((b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}(G_{43}(s_{(40)}), s_{(40)})\right) T_{41}(s_{(40)})\right] ds_{(40)}$ | |

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| $\overline{T}_{42}(t) = T_{42}^0 + \int_0^t \left[(b_{42})^{(8)} T_{41}(s_{(40)}) - \left((b'_{42})^{(8)} - (b''_{42})^{(8)}(G_{43}(s_{(40)}), s_{(40)})\right) T_{42}(s_{(40)})\right] ds_{(40)}$<br><br>Where $s_{(40)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(9)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 166<br>A |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{44})^{(9)}$, $T_i^0 \le (\hat{Q}_{44})^{(9)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{44})^{(9)} e^{(\bar{M}_{44})^{(9)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{44})^{(9)} e^{(\bar{M}_{44})^{(9)} t}$ | |
| By<br><br>$\bar{G}_{44}(t) = G_{44}^0 + \int_0^t \left[(a_{44})^{(9)} G_{45}(s_{(44)}) - \left((a'_{44})^{(9)} + a''_{44})^{(9)}(T_{45}(s_{(44)}), s_{(44)})\right) G_{44}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{G}_{45}(t) = G_{45}^0 + \int_0^t \left[(a_{45})^{(9)} G_{44}(s_{(44)}) - \left((a'_{45})^{(9)} + (a''_{45})^{(9)}(T_{45}(s_{(44)}), s_{(44)})\right) G_{45}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{G}_{46}(t) = G_{46}^0 + \int_0^t \left[(a_{46})^{(9)} G_{45}(s_{(44)}) - \left((a'_{46})^{(9)} + (a''_{46})^{(9)}(T_{45}(s_{(44)}), s_{(44)})\right) G_{46}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{44}(t) = T_{44}^0 + \int_0^t \left[(b_{44})^{(9)} T_{45}(s_{(44)}) - \left((b'_{44})^{(9)} - (b''_{44})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{44}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{45}(t) = T_{45}^0 + \int_0^t \left[(b_{45})^{(9)} T_{44}(s_{(44)}) - \left((b'_{45})^{(9)} - (b''_{45})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{45}(s_{(44)})\right] ds_{(44)}$ | |
| $\overline{T}_{46}(t) = T_{46}^0 + \int_0^t \left[(b_{46})^{(9)} T_{45}(s_{(44)}) - \left((b'_{46})^{(9)} - (b''_{46})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{46}(s_{(44)})\right] ds_{(44)}$<br><br>Where $s_{(44)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| The operator $\mathcal{A}^{(1)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{13}(t) \le G_{13}^0 + \int_0^t \left[(a_{13})^{(1)} \left(G_{14}^0 + (\hat{P}_{13})^{(1)} e^{(\bar{M}_{13})^{(1)} s_{(13)}}\right)\right] ds_{(13)} =$<br><br>$\left(1 + (a_{13})^{(1)} t\right) G_{14}^0 + \frac{(a_{13})^{(1)} (\hat{P}_{13})^{(1)}}{(\bar{M}_{13})^{(1)}} \left(e^{(\bar{M}_{13})^{(1)} t} - 1\right)$ | 167 |
| From which it follows that<br><br>$(G_{13}(t) - G_{13}^0) e^{-(\bar{M}_{13})^{(1)} t} \le \frac{(a_{13})^{(1)}}{(\bar{M}_{13})^{(1)}} \left[\left((\hat{P}_{13})^{(1)} + G_{14}^0\right) e^{\left(-\frac{(\hat{P}_{13})^{(1)} + G_{14}^0}{G_{14}^0}\right)} + (\hat{P}_{13})^{(1)}\right]$<br><br>$(G_i^0)$ is as defined in the statement of theorem 1 | 168 |
| Analogous inequalities hold also for $G_{14}, G_{15}, T_{13}, T_{14}, T_{15}$ | |

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| The operator $\mathcal{A}^{(2)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that | |
| $G_{16}(t) \le G_{16}^0 + \int_0^t \left[(a_{16})^{(2)}\left(G_{17}^0 + (\hat{P}_{16})^{(6)} e^{(\hat{M}_{16})^{(2)} s_{(16)}}\right)\right] ds_{(16)} =$ $\left(1 + (a_{16})^{(2)} t\right) G_{17}^0 + \frac{(a_{16})^{(2)} (\hat{P}_{16})^{(2)}}{(\hat{M}_{16})^{(2)}}\left(e^{(\hat{M}_{16})^{(2)} t} - 1\right)$ | 169 |
| From which it follows that $(G_{16}(t) - G_{16}^0) e^{-(\hat{M}_{16})^{(2)} t} \le \frac{(a_{16})^{(2)}}{(\hat{M}_{16})^{(2)}}\left[\left((\hat{P}_{16})^{(2)} + G_{17}^0\right) e^{\left(-\frac{(\hat{P}_{16})^{(2)} + G_{17}^0}{G_{17}^0}\right)} + (\hat{P}_{16})^{(2)}\right]$ | 170 |
| Analogous inequalities hold also for $G_{17}, G_{18}, T_{16}, T_{17}, T_{18}$ | |
| The operator $\mathcal{A}^{(3)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{20}(t) \le G_{20}^0 + \int_0^t \left[(a_{20})^{(3)}\left(G_{21}^0 + (\hat{P}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} s_{(20)}}\right)\right] ds_{(20)} =$ $\left(1 + (a_{20})^{(3)} t\right) G_{21}^0 + \frac{(a_{20})^{(3)} (\hat{P}_{20})^{(3)}}{(\hat{M}_{20})^{(3)}}\left(e^{(\hat{M}_{20})^{(3)} t} - 1\right)$ | 171 |
| From which it follows that $(G_{20}(t) - G_{20}^0) e^{-(\hat{M}_{20})^{(3)} t} \le \frac{(a_{20})^{(3)}}{(\hat{M}_{20})^{(3)}}\left[\left((\hat{P}_{20})^{(3)} + G_{21}^0\right) e^{\left(-\frac{(\hat{P}_{20})^{(3)} + G_{21}^0}{G_{21}^0}\right)} + (\hat{P}_{20})^{(3)}\right]$ | 172 |
| Analogous inequalities hold also for $G_{21}, G_{22}, T_{20}, T_{21}, T_{22}$ | |
| The operator $\mathcal{A}^{(4)}$ maps the space of functions satisfying into itself .Indeed it is obvious that $G_{24}(t) \le G_{24}^0 + \int_0^t \left[(a_{24})^{(4)}\left(G_{25}^0 + (\hat{P}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} s_{(24)}}\right)\right] ds_{(24)} =$ $\left(1 + (a_{24})^{(4)} t\right) G_{25}^0 + \frac{(a_{24})^{(4)} (\hat{P}_{24})^{(4)}}{(\hat{M}_{24})^{(4)}}\left(e^{(\hat{M}_{24})^{(4)} t} - 1\right)$ | 173 |
| From which it follows that $(G_{24}(t) - G_{24}^0) e^{-(\hat{M}_{24})^{(4)} t} \le \frac{(a_{24})^{(4)}}{(\hat{M}_{24})^{(4)}}\left[\left((\hat{P}_{24})^{(4)} + G_{25}^0\right) e^{\left(-\frac{(\hat{P}_{24})^{(4)} + G_{25}^0}{G_{25}^0}\right)} + (\hat{P}_{24})^{(4)}\right]$ $(G_i^0)$ is as defined in the statement of theorem 4 | 174 |
| The operator $\mathcal{A}^{(5)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{28}(t) \le G_{28}^0 + \int_0^t \left[(a_{28})^{(5)}\left(G_{29}^0 + (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} s_{(28)}}\right)\right] ds_{(28)} =$ $\left(1 + (a_{28})^{(5)} t\right) G_{29}^0 + \frac{(a_{28})^{(5)} (\hat{P}_{28})^{(5)}}{(\hat{M}_{28})^{(5)}}\left(e^{(\hat{M}_{28})^{(5)} t} - 1\right)$ | |

From which it follows that

$$(G_{28}(t)-G_{28}^{0})e^{-(\hat{M}_{28})^{(5)}t} \leq \frac{(a_{28})^{(5)}}{(\hat{M}_{28})^{(5)}}\left[((\hat{P}_{28})^{(5)}+G_{29}^{0})e^{\left(-\frac{(\hat{P}_{28})^{(5)}+G_{29}^{0}}{G_{29}^{0}}\right)}+(\hat{P}_{28})^{(5)}\right]$$

$(G_i^0)$ is as defined in the statement of theorem 5 — 175

The operator $\mathcal{A}^{(6)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that — 176

$$G_{32}(t) \leq G_{32}^{0}+\int_0^t\left[(a_{32})^{(6)}\left(G_{33}^{0}+(\hat{P}_{32})^{(6)}e^{(\hat{M}_{32})^{(6)}s_{(32)}}\right)\right]ds_{(32)} =$$

$$\left(1+(a_{32})^{(6)}t\right)G_{33}^{0}+\frac{(a_{32})^{(6)}(\hat{P}_{32})^{(6)}}{(\hat{M}_{32})^{(6)}}\left(e^{(\hat{M}_{32})^{(6)}t}-1\right)$$

From which it follows that — 177

$$(G_{32}(t)-G_{32}^{0})e^{-(\hat{M}_{32})^{(6)}t} \leq \frac{(a_{32})^{(6)}}{(\hat{M}_{32})^{(6)}}\left[((\hat{P}_{32})^{(6)}+G_{33}^{0})e^{\left(-\frac{(\hat{P}_{32})^{(6)}+G_{33}^{0}}{G_{33}^{0}}\right)}+(\hat{P}_{32})^{(6)}\right]$$

$(G_i^0)$ is as defined in the statement of theorem 6

Analogous inequalities hold also for $G_{25}, G_{26}, T_{24}, T_{25}, T_{26}$

(e) The operator $\mathcal{A}^{(7)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that — 178

$$G_{36}(t) \leq G_{36}^{0}+\int_0^t\left[(a_{36})^{(7)}\left(G_{37}^{0}+(\hat{P}_{36})^{(7)}e^{(\hat{M}_{36})^{(7)}s_{(36)}}\right)\right]ds_{(36)} =$$

$$\left(1+(a_{36})^{(7)}t\right)G_{37}^{0}+\frac{(a_{36})^{(7)}(\hat{P}_{36})^{(7)}}{(\hat{M}_{36})^{(7)}}\left(e^{(\hat{M}_{36})^{(7)}t}-1\right)$$

From which it follows that

$$(G_{36}(t)-G_{36}^{0})e^{-(\hat{M}_{36})^{(7)}t} \leq \frac{(a_{36})^{(7)}}{(\hat{M}_{36})^{(7)}}\left[((\hat{P}_{36})^{(7)}+G_{37}^{0})e^{\left(-\frac{(\hat{P}_{36})^{(7)}+G_{37}^{0}}{G_{37}^{0}}\right)}+(\hat{P}_{36})^{(7)}\right]$$

$(G_i^0)$ is as defined in the statement of theorem 7

The operator $\mathcal{A}^{(8)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that

$$G_{40}(t) \leq G_{40}^{0}+\int_0^t\left[(a_{40})^{(8)}\left(G_{41}^{0}+(\hat{P}_{40})^{(8)}e^{(\hat{M}_{40})^{(8)}s_{(40)}}\right)\right]ds_{(40)} =$$

$$\left(1+(a_{40})^{(8)}t\right)G_{41}^{0}+\frac{(a_{40})^{(8)}(\hat{P}_{40})^{(8)}}{(\hat{M}_{40})^{(8)}}\left(e^{(\hat{M}_{40})^{(8)}t}-1\right)$$ — 180

From which it follows that

$$(G_{40}(t)-G_{40}^{0})e^{-(\hat{M}_{40})^{(8)}t} \leq \frac{(a_{40})^{(8)}}{(\hat{M}_{40})^{(8)}}\left[((\hat{P}_{40})^{(8)}+G_{41}^{0})e^{\left(-\frac{(\hat{P}_{40})^{(8)}+G_{41}^{0}}{G_{41}^{0}}\right)}+(\hat{P}_{40})^{(8)}\right]$$ — 181

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| $(G_i^0)$ is as defined in the statement of theorem 8<br>Analogous inequalities hold also for $G_{41}, G_{42}, T_{40}, T_{41}, T_{42}$ | |
| The operator $\mathcal{A}^{(9)}$ maps the space of functions satisfying 34,35,36 into itself .Indeed it is obvious that<br>$G_{44}(t) \le G_{44}^0 + \int_0^t \left[(a_{44})^{(9)}\left(G_{45}^0 + (\hat{P}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} s_{(44)}}\right)\right] ds_{(44)} =$<br>$\left(1 + (a_{44})^{(9)} t\right) G_{45}^0 + \frac{(a_{44})^{(9)} (\hat{P}_{44})^{(9)}}{(\hat{M}_{44})^{(9)}}\left(e^{(\hat{M}_{44})^{(9)} t} - 1\right)$ | |
| From which it follows that<br>$(G_{44}(t) - G_{44}^0) e^{-(\hat{M}_{44})^{(9)} t} \le \frac{(a_{44})^{(9)}}{(\hat{M}_{44})^{(9)}}\left[\left((\hat{P}_{44})^{(9)} + G_{45}^0\right) e^{\left(-\frac{(\hat{P}_{44})^{(9)} + G_{45}^0}{G_{45}^0}\right)} + (\hat{P}_{44})^{(9)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 9<br>Analogous inequalities hold also for $G_{45}, G_{46}, T_{44}, T_{45}, T_{46}$ | |
| It is now sufficient to take $\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}} < 1$ and to choose<br>$(\hat{P}_{13})^{(1)}$ and $(\hat{Q}_{13})^{(1)}$ large to have | 182 |
| $\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}\left[(\hat{P}_{13})^{(1)} + \left((\hat{P}_{13})^{(1)} + G_j^0\right) e^{-\left(\frac{(\hat{P}_{13})^{(1)} + G_j^0}{G_j^0}\right)}\right] \le (\hat{P}_{13})^{(1)}$ | 183 |
| $\frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}}\left[\left((\hat{Q}_{13})^{(1)} + T_j^0\right) e^{-\left(\frac{(\hat{Q}_{13})^{(1)} + T_j^0}{T_j^0}\right)} + (\hat{Q}_{13})^{(1)}\right] \le (\hat{Q}_{13})^{(1)}$ | 184 |
| In order that the operator $\mathcal{A}^{(1)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(1)}$ is a contraction with respect to the metric<br>$d\left(\left(G^{(1)}, T^{(1)}\right), \left(G^{(2)}, T^{(2)}\right)\right) =$<br>$\sup_i \{\max_{t \in \mathbb{R}_+} \left|G_i^{(1)}(t) - G_i^{(2)}(t)\right| e^{-(\hat{M}_{13})^{(1)} t}, \max_{t \in \mathbb{R}_+} \left|T_i^{(1)}(t) - T_i^{(2)}(t)\right| e^{-(\hat{M}_{13})^{(1)} t}\}$ | 185 |
| Indeed if we denote<br>**Definition of** $\tilde{G}, \tilde{T}$ : $\left(\tilde{G}, \tilde{T}\right) = \mathcal{A}^{(1)}(G, T)$<br>It results | |

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| $\left\|\tilde{G}_{13}^{(1)}-\tilde{G}_{i}^{(2)}\right\| \leq \int_0^t (a_{13})^{(1)} \left\|G_{14}^{(1)}-G_{14}^{(2)}\right\| e^{-(\widehat{M}_{13})^{(1)}s_{(13)}} e^{(\widehat{M}_{13})^{(1)}s_{(13)}}\, ds_{(13)} +$<br>$\int_0^t \{(a_{13}')^{(1)} \left\|G_{13}^{(1)}-G_{13}^{(2)}\right\| e^{-(\widehat{M}_{13})^{(1)}s_{(13)}} e^{-(\widehat{M}_{13})^{(1)}s_{(13)}} +$<br>$(a_{13}'')^{(1)}\left(T_{14}^{(1)}, s_{(13)}\right)\left\|G_{13}^{(1)}-G_{13}^{(2)}\right\| e^{-(\widehat{M}_{13})^{(1)}s_{(13)}} e^{(\widehat{M}_{13})^{(1)}s_{(13)}} +$<br>$G_{13}^{(2)} \left\|(a_{13}'')^{(1)}\left(T_{14}^{(1)}, s_{(13)}\right) - (a_{13}'')^{(1)}\left(T_{14}^{(2)}, s_{(13)}\right)\right\| \; e^{-(\widehat{M}_{13})^{(1)}s_{(13)}} e^{(\widehat{M}_{13})^{(1)}s_{(13)}}\} ds_{(13)}$<br>Where $s_{(13)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | |
| $\left\|G^{(1)}-G^{(2)}\right\| e^{-(\widehat{M}_{13})^{(1)}t} \leq$<br>$\frac{1}{(\widehat{M}_{13})^{(1)}}\left((a_{13})^{(1)} + (a_{13}')^{(1)} + (\widehat{A}_{13})^{(1)} + (\widehat{P}_{13})^{(1)}(\widehat{k}_{13})^{(1)}\right) d\left(\left(G^{(1)}, T^{(1)}; G^{(2)}, T^{(2)}\right)\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 186 |
| **Remark 1:** The fact that we supposed $(a_{13}'')^{(1)}$ and $(b_{13}'')^{(1)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{13})^{(1)} e^{(\widehat{M}_{13})^{(1)}t}$ *and* $(\widehat{Q}_{13})^{(1)} e^{(\widehat{M}_{13})^{(1)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i'')^{(1)}$ and $(b_i'')^{(1)}, i = 13,14,15$ depend only on $T_{14}$ and respectively on $G$(*and not on* $t$) and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 2:** There does not exist any $t$ where $G_i(t) = 0$ *and* $T_i(t) = 0$<br>From 19 to 24 it results<br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i')^{(1)} - (a_i'')^{(1)}(T_{14}(s_{(13)}), s_{(13)})\} ds_{(13)}\right]} \geq 0$<br>$T_i(t) \geq T_i^0 e^{\left(-(b_i')^{(1)}t\right)} > 0$ for t > 0 | |
| **Definition of** $\left((\widehat{M}_{13})^{(1)}\right)_1, \left((\widehat{M}_{13})^{(1)}\right)_2$ *and* $\left((\widehat{M}_{13})^{(1)}\right)_3$ :<br>**Remark 3**: if $G_{13}$ is bounded, the same property have also $G_{14}$ *and* $G_{15}$ . indeed if<br>$G_{13} < (\widehat{M}_{13})^{(1)}$ it follows $\frac{dG_{14}}{dt} \leq \left((\widehat{M}_{13})^{(1)}\right)_1 - (a_{14}')^{(1)} G_{14}$ and by integrating<br>$G_{14} \leq \left((\widehat{M}_{13})^{(1)}\right)_2 = G_{14}^0 + 2(a_{14})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_1 / (a_{14}')^{(1)}$<br>In the same way , one can obtain<br>$G_{15} \leq \left((\widehat{M}_{13})^{(1)}\right)_3 = G_{15}^0 + 2(a_{15})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_2 / (a_{15}')^{(1)}$ | 187 |

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| If $G_{14}$ *or* $G_{15}$ is bounded, the same property follows for $G_{13}$ , $G_{15}$ and $G_{13}$ , $G_{14}$ respectively. | |
| **Remark 4:** If $G_{13}$ *is* bounded, from below, the same property holds for $G_{14}$ *and* $G_{15}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{14}$ is bounded from below. | 188 |
| **Remark 5:** If $T_{13}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(1)}(G(t),t)) = (b_{14}^{'})^{(1)}$ then $T_{14} \to \infty$.<br><br>**Definition of** $(m)^{(1)}$ and $\varepsilon_1$ :<br><br>Indeed let $t_1$ be so that for $t > t_1$<br><br>$(b_{14})^{(1)} - (b_i^{''})^{(1)}(G(t),t) < \varepsilon_1, T_{13}(t) > (m)^{(1)}$ | 189 |
| Then $\frac{dT_{14}}{dt} \geq (a_{14})^{(1)}(m)^{(1)} - \varepsilon_1 T_{14}$ which leads to<br><br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{\varepsilon_1}\right)(1-e^{-\varepsilon_1 t}) + T_{14}^0 e^{-\varepsilon_1 t}$ If we take t such that $e^{-\varepsilon_1 t} = \frac{1}{2}$ it results<br><br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{2}\right)$, $t = log\frac{2}{\varepsilon_1}$ By taking now $\varepsilon_1$ sufficiently small one sees that $T_{14}$ is unbounded.<br>The same property holds for $T_{15}$ if $\lim_{t\to\infty}(b_{15}^{''})^{(1)}(G(t),t) = (b_{15}^{'})^{(1)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |
| It is now sufficient to take $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}$ , $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ and to choose<br><br>$(\hat{P}_{16})^{(2)}$ *and* $(\hat{Q}_{16})^{(2)}$ large to have | 190 |
| $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[(\hat{P}_{16})^{(2)} + \left((\hat{P}_{16})^{(2)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{16})^{(2)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{16})^{(2)}$ | 191 |
| $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[\left((\hat{Q}_{16})^{(2)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{16})^{(2)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{16})^{(2)}\right] \leq (\hat{Q}_{16})^{(2)}$ | 192 |
| In order that the operator $\mathcal{A}^{(2)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 193 |
| The operator $\mathcal{A}^{(2)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{19})^{(1)},(T_{19})^{(1)}),((G_{19})^{(2)},(T_{19})^{(2)})\right) =$<br><br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t) - G_i^{(2)}(t)\right|e^{-(\hat{M}_{16})^{(2)}t}, \max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t) - T_i^{(2)}(t)\right|e^{-(\hat{M}_{16})^{(2)}t}\}$ | 194 |
| Indeed if we denote | 195 |

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| **Definition of** $\widetilde{G_{19}}, \widetilde{T_{19}}$ : $\left(\widetilde{G_{19}}, \widetilde{T_{19}}\right) = \mathcal{A}^{(2)}(G_{19}, T_{19})$ | |
| It results<br>$\left\|\tilde{G}_{16}^{(1)} - \tilde{G}_{i}^{(2)}\right\| \leq \int_0^t (a_{16})^{(2)} \left\|G_{17}^{(1)} - G_{17}^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{(\widehat{M}_{16})^{(2)} s_{(16)}}\, ds_{(16)} +$<br>$\int_0^t \{(a_{16}^{'})^{(2)} \left\|G_{16}^{(1)} - G_{16}^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} +$<br>$(a_{16}^{''})^{(2)}\left(T_{17}^{(1)}, s_{(16)}\right)\left\|G_{16}^{(1)} - G_{16}^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{(\widehat{M}_{16})^{(2)} s_{(16)}} +$<br>$G_{16}^{(2)} \left\|(a_{16}^{''})^{(2)}\left(T_{17}^{(1)}, s_{(16)}\right) - (a_{16}^{''})^{(2)}\left(T_{17}^{(2)}, s_{(16)}\right)\right\| \; e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{(\widehat{M}_{16})^{(2)} s_{(16)}} \} ds_{(16)}$ | 196 |
| Where $s_{(16)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | 197 |
| $\left\|(G_{19})^{(1)} - (G_{19})^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} t} \leq$<br>$\frac{1}{(\widehat{M}_{16})^{(2)}}\left((a_{16})^{(2)} + (a_{16}^{'})^{(2)} + (\widehat{A}_{16})^{(2)} + (\widehat{P}_{16})^{(2)}(\widehat{k}_{16})^{(2)}\right) d\left(((G_{19})^{(1)}, (T_{19})^{(1)}; (G_{19})^{(2)}, (T_{19})^{(2)})\right)$ | |
| And analogous inequalities for $G_i$ and $T_i$. Taking into account the hypothesis the result follows | 198 |
| **Remark 6:** The fact that we supposed $(a_{16}^{''})^{(2)}$ and $(b_{16}^{''})^{(2)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{16})^{(2)} e^{(\widehat{M}_{16})^{(2)} t}$ and $(\widehat{Q}_{16})^{(2)} e^{(\widehat{M}_{16})^{(2)} t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}, i = 16,17,18$ depend only on $T_{17}$ and respectively on $(G_{19})$(and not on t) and hypothesis can replaced by a usual Lipschitz condition. | 199 |
| **Remark 7:** There does not exist any t where $G_i(t) = 0$ and $T_i(t) = 0$<br>it results<br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(2)} - (a_i^{''})^{(2)}(T_{17}(s_{(16)}), s_{(16)})\} ds_{(16)}\right]} \geq 0$<br>$T_i(t) \geq T_i^0 e^{(-(b_i^{'})^{(2)} t)} > 0$ for $t > 0$ | 200 |
| **Definition of** $\left((\widehat{M}_{16})^{(2)}\right)_1$, $\left((\widehat{M}_{16})^{(2)}\right)_2$ and $\left((\widehat{M}_{16})^{(2)}\right)_3$ :<br>**Remark 8:** if $G_{16}$ is bounded, the same property have also $G_{17}$ and $G_{18}$ . indeed if<br>$G_{16} < (\widehat{M}_{16})^{(2)}$ it follows $\frac{dG_{17}}{dt} \leq \left((\widehat{M}_{16})^{(2)}\right)_1 - (a_{17}^{'})^{(2)} G_{17}$ and by integrating<br>$G_{17} \leq \left((\widehat{M}_{16})^{(2)}\right)_2 = G_{17}^0 + 2(a_{17})^{(2)}\left((\widehat{M}_{16})^{(2)}\right)_1 / (a_{17}^{'})^{(2)}$ | 201 |

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| In the same way , one can obtain<br>$G_{18} \leq \left((\widehat{M}_{16})^{(2)}\right)_3 = G_{18}^0 + 2(a_{18})^{(2)}\left((\widehat{M}_{16})^{(2)}\right)_2/(a_{18}')^{(2)}$<br>If $G_{17}$ or $G_{18}$ is bounded, the same property follows for $G_{16}$ , $G_{18}$ and $G_{16}$ , $G_{17}$ respectively. | |
| **Remark 9:** If $G_{16}$ is bounded, from below, the same property holds for $G_{17}$ and $G_{18}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{17}$ is bounded from below. | 202 |
| **Remark 10:** If $T_{16}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(2)}((G_{19})(t),t)) = (b_{17}')^{(2)}$ then $T_{17} \to \infty$.<br>**Definition of** $(m)^{(2)}$ and $\varepsilon_2$ :<br>Indeed let $t_2$ be so that for $t > t_2$<br>$(b_{17})^{(2)} - (b_i'')^{(2)}((G_{19})(t),t) < \varepsilon_2, T_{16}(t) > (m)^{(2)}$ | 203 |
| Then $\frac{dT_{17}}{dt} \geq (a_{17})^{(2)}(m)^{(2)} - \varepsilon_2 T_{17}$ which leads to<br>$T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{\varepsilon_2}\right)(1 - e^{-\varepsilon_2 t}) + T_{17}^0 e^{-\varepsilon_2 t}$ If we take t such that $e^{-\varepsilon_2 t} = \frac{1}{2}$ it results | 204 |
| $T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{2}\right)$, $t = \log\frac{2}{\varepsilon_2}$ By taking now $\varepsilon_2$ sufficiently small one sees that $T_{17}$ is unbounded.<br>The same property holds for $T_{18}$ if $\lim_{t\to\infty}(b_{18}'')^{(2)}\left((G_{19})(t),t\right) = (b_{18}')^{(2)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 205 |
| It is now sufficient to take $\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}$ , $\frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}} < 1$ and to choose<br>$(\hat{P}_{20})^{(3)}$ and $(\hat{Q}_{20})^{(3)}$ large to have | 207 |
| $\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}\left[(\hat{P}_{20})^{(3)} + \left((\hat{P}_{20})^{(3)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{20})^{(3)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{20})^{(3)}$ | 208 |
| $\frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}}\left[\left((\hat{Q}_{20})^{(3)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{20})^{(3)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{20})^{(3)}\right] \leq (\hat{Q}_{20})^{(3)}$ | 209 |
| In order that the operator $\mathcal{A}^{(3)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 210 |
| The operator $\mathcal{A}^{(3)}$ is a contraction with respect to the metric<br>$d\left(((G_{23})^{(1)},(T_{23})^{(1)}),((G_{23})^{(2)},(T_{23})^{(2)})\right) =$ | 211 |

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| $\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{20})^{(3)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{20})^{(3)}t}\}$ | |
| Indeed if we denote<br>**Definition of** $\widetilde{G_{23}},\widetilde{T_{23}}:(\widetilde{(G_{23})},\widetilde{(T_{23})})=\mathcal{A}^{(3)}((G_{23}),(T_{23}))$ | 212 |
| It results<br>$\left\|\tilde{G}_{20}^{(1)}-\tilde{G}_i^{(2)}\right\|\le\int_0^t(a_{20})^{(3)}\left\|G_{21}^{(1)}-G_{21}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{(\widehat{M}_{20})^{(3)}s_{(20)}}ds_{(20)}+$<br>$\int_0^t\{(a_{20}^{'})^{(3)}\left\|G_{20}^{(1)}-G_{20}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}+$<br>$(a_{20}^{''})^{(3)}\left(T_{21}^{(1)},s_{(20)}\right)\left\|G_{20}^{(1)}-G_{20}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{(\widehat{M}_{20})^{(3)}s_{(20)}}+$<br>$G_{20}^{(2)}\left\|(a_{20}^{''})^{(3)}\left(T_{21}^{(1)},s_{(20)}\right)-(a_{20}^{''})^{(3)}\left(T_{21}^{(2)},s_{(20)}\right)\right\|\ e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{(\widehat{M}_{20})^{(3)}s_{(20)}}\}ds_{(20)}$<br>Where $s_{(20)}$ represents integrand that is integrated over the interval $[0,t]$<br>From the hypotheses it follows | 213 |
| $\left\|G_{23}{}^{(1)}-G_{23}{}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}t}\le$<br>$\frac{1}{(\widehat{M}_{20})^{(3)}}\left((a_{20})^{(3)}+(a_{20}^{'})^{(3)}+(\widehat{A}_{20})^{(3)}+(\widehat{P}_{20})^{(3)}(\widehat{k}_{20})^{(3)}\right)d\left(((G_{23})^{(1)},(T_{23})^{(1)};\ (G_{23})^{(2)},(T_{23})^{(2)})\right)$<br>And analogous inequalities for $G_i\ and\ T_i$. Taking into account the hypothesis the result follows | 214 |
| **Remark 11:** The fact that we supposed $(a_{20}^{''})^{(3)}$ and $(b_{20}^{''})^{(3)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}\ and\ (\widehat{Q}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(3)}$ and $(b_i^{''})^{(3)}, i=20,21,22$ depend only on $T_{21}$ and respectively on $(G_{23})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 215 |
| **Remark 12:** There does not exist any $t$ where $G_i(t)=0\ and\ T_i(t)=0$<br>it results<br>$G_i(t)\ge G_i^0e^{\left[-\int_0^t\{(a_i^{'})^{(3)}-(a_i^{''})^{(3)}(T_{21}(s_{(20)}),s_{(20)})\}ds_{(20)}\right]}\ge 0$<br>$T_i(t)\ge T_i^0e^{(-(b_i^{'})^{(3)}t)}>0$ for t > 0 | 216 |

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| **Definition of** $\left((\widehat{M}_{20})^{(3)}\right)_1, \left((\widehat{M}_{20})^{(3)}\right)_2$ *and* $\left((\widehat{M}_{20})^{(3)}\right)_3$ :<br><br>**Remark 13:** if $G_{20}$ is bounded, the same property have also $G_{21}$ *and* $G_{22}$ . indeed if<br><br>$G_{20} < (\widehat{M}_{20})^{(3)}$ it follows $\frac{dG_{21}}{dt} \leq \left((\widehat{M}_{20})^{(3)}\right)_1 - (a_{21}^{'})^{(3)} G_{21}$ and by integrating<br><br>$G_{21} \leq \left((\widehat{M}_{20})^{(3)}\right)_2 = G_{21}^0 + 2(a_{21})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_1/(a_{21}^{'})^{(3)}$<br><br>In the same way , one can obtain<br><br>$G_{22} \leq \left((\widehat{M}_{20})^{(3)}\right)_3 = G_{22}^0 + 2(a_{22})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_2/(a_{22}^{'})^{(3)}$<br><br>If $G_{21}$ *or* $G_{22}$ is bounded, the same property follows for $G_{20}$ , $G_{22}$ and $G_{20}$ , $G_{21}$ respectively. | 217 |
| **Remark 14:** If $G_{20}$ *is* bounded, from below, the same property holds for $G_{21}$ *and* $G_{22}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{21}$ is bounded from below. | 218 |
| **Remark 15:** If $T_{20}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(3)}\left((G_{23})(t),t\right) = (b_{21}^{'})^{(3)}$ then $T_{21} \to \infty$.<br><br>**Definition of** $(m)^{(3)}$ and $\varepsilon_3$ :<br><br>Indeed let $t_3$ be so that for $t > t_3$<br><br>$(b_{21})^{(3)} - (b_i^{''})^{(3)}\left((G_{23})(t),t\right) < \varepsilon_3, T_{20}(t) > (m)^{(3)}$ | 219 |
| Then $\frac{dT_{21}}{dt} \geq (a_{21})^{(3)}(m)^{(3)} - \varepsilon_3 T_{21}$ which leads to<br><br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{\varepsilon_3}\right)(1 - e^{-\varepsilon_3 t}) + T_{21}^0 e^{-\varepsilon_3 t}$ If we take t such that $e^{-\varepsilon_3 t} = \frac{1}{2}$ it results<br><br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{2}\right)$, $t = log\frac{2}{\varepsilon_3}$ By taking now $\varepsilon_3$ sufficiently small one sees that $T_{21}$ is unbounded.<br>The same property holds for $T_{22}$ if $\lim_{t\to\infty}(b_{22}^{''})^{(3)}\left((G_{23})(t),t\right) = (b_{22}^{'})^{(3)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 220 |
| It is now sufficient to take $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}$ , $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}} < 1$ and to choose<br><br>$(\widehat{P}_{24})^{(4)}$ and $(\widehat{Q}_{24})^{(4)}$ large to have | 221 |
| $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[(\widehat{P}_{24})^{(4)} + \left((\widehat{P}_{24})^{(4)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{24})^{(4)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{24})^{(4)}$ | 222 |
| $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[\left((\widehat{Q}_{24})^{(4)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{24})^{(4)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{24})^{(4)}\right] \leq (\widehat{Q}_{24})^{(4)}$ | 223 |

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| In order that the operator $\mathcal{A}^{(4)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | 224 |
| The operator $\mathcal{A}^{(4)}$ is a contraction with respect to the metric<br>$d\left(((G_{27})^{(1)},(T_{27})^{(1)}),((G_{27})^{(2)},(T_{27})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{24})^{(4)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{24})^{(4)}t}\}$<br>Indeed if we denote<br>**Definition of** $(\widetilde{G_{27}}),(\widetilde{T_{27}})$ : $\left((\widetilde{G_{27}}),(\widetilde{T_{27}})\right)=\mathcal{A}^{(4)}((G_{27}),(T_{27}))$<br>It results<br>$\left\|\tilde{G}_{24}^{(1)}-\tilde{G}_i^{(2)}\right\|\le\int_0^t(a_{24})^{(4)}\left\|G_{25}^{(1)}-G_{25}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}\,ds_{(24)}+$<br>$\int_0^t\{(a_{24}^{'})^{(4)}\left\|G_{24}^{(1)}-G_{24}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}+$<br>$(a_{24}^{''})^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)\left\|G_{24}^{(1)}-G_{24}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}+$<br>$G_{24}^{(2)}\left\|(a_{24}^{''})^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)-(a_{24}^{''})^{(4)}\left(T_{25}^{(2)},s_{(24)}\right)\right\|\ e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}\}ds_{(24)}$<br>Where $s_{(24)}$ represents integrand that is integrated over the interval $[0,\mathrm{t}]$<br>From the hypotheses on Equations it follows | 225 |
| $\left\|(G_{27})^{(1)}-(G_{27})^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}t}\le$<br>$\frac{1}{(\widehat{M}_{24})^{(4)}}\left((a_{24})^{(4)}+(a_{24}^{'})^{(4)}+(\widehat{A}_{24})^{(4)}+(\widehat{P}_{24})^{(4)}(\widehat{k}_{24})^{(4)}\right)d\left(((G_{27})^{(1)},(T_{27})^{(1)};\ (G_{27})^{(2)},(T_{27})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 226 |
| **Remark 16:** The fact that we supposed $(a_{24}^{''})^{(4)}$ and $(b_{24}^{''})^{(4)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}$ *and* $(\widehat{Q}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(4)}$ and $(b_i^{''})^{(4)}, i=24{,}25{,}26$ depend only on $\mathrm{T}_{25}$ and respectively on $(G_{27})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 227 |
| **Remark 17:** There does not exist any $t$ where $G_i(t)=0$ *and* $T_i(t)=0$<br>it results | 228 |

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| $G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i')^{(4)} - (a_i'')^{(4)}(T_{25}(s_{(24)}), s_{(24)})\} ds_{(24)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{\left(-(b_i')^{(4)} t\right)} > 0$ for t > 0 | |
| **Definition of** $\left((\widehat{M}_{24})^{(4)}\right)_1, \left((\widehat{M}_{24})^{(4)}\right)_2$ *and* $\left((\widehat{M}_{24})^{(4)}\right)_3$ :<br><br>**Remark 18:** if $G_{24}$ is bounded, the same property have also $G_{25}$ *and* $G_{26}$ . indeed if<br><br>$G_{24} < (\widehat{M}_{24})^{(4)}$ it follows $\frac{dG_{25}}{dt} \leq \left((\widehat{M}_{24})^{(4)}\right)_1 - (a_{25}')^{(4)} G_{25}$ and by integrating<br><br>$G_{25} \leq \left((\widehat{M}_{24})^{(4)}\right)_2 = G_{25}^0 + 2(a_{25})^{(4)} \left((\widehat{M}_{24})^{(4)}\right)_1 / (a_{25}')^{(4)}$<br><br>In the same way , one can obtain<br><br>$G_{26} \leq \left((\widehat{M}_{24})^{(4)}\right)_3 = G_{26}^0 + 2(a_{26})^{(4)} \left((\widehat{M}_{24})^{(4)}\right)_2 / (a_{26}')^{(4)}$<br><br>If $G_{25}$ *or* $G_{26}$ is bounded, the same property follows for $G_{24}$ , $G_{26}$ and $G_{24}$ , $G_{25}$ respectively. | 229 |
| **Remark 19:** If $G_{24}$ *is* bounded, from below, the same property holds for $G_{25}$ *and* $G_{26}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{25}$ is bounded from below. | 230 |
| **Remark 20:** If $T_{24}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(4)}((G_{27})(t), t)) = (b_{25}')^{(4)}$ then $T_{25} \to \infty$.<br><br>**Definition of** $(m)^{(4)}$ and $\varepsilon_4$ :<br><br>Indeed let $t_4$ be so that for $t > t_4$<br><br>$(b_{25})^{(4)} - (b_i'')^{(4)}((G_{27})(t), t) < \varepsilon_4, T_{24}(t) > (m)^{(4)}$ | 231 |
| Then $\frac{dT_{25}}{dt} \geq (a_{25})^{(4)}(m)^{(4)} - \varepsilon_4 T_{25}$ which leads to<br><br>$T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{\varepsilon_4}\right)(1 - e^{-\varepsilon_4 t}) + T_{25}^0 e^{-\varepsilon_4 t}$ If we take t such that $e^{-\varepsilon_4 t} = \frac{1}{2}$ it results<br><br>$T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{2}\right)$, $t = log\frac{2}{\varepsilon_4}$ By taking now $\varepsilon_4$ sufficiently small one sees that $T_{25}$ is unbounded. The same property holds for $T_{26}$ if $\lim_{t\to\infty}(b_{26}'')^{(4)}\left((G_{27})(t), t\right) = (b_{26}')^{(4)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 42<br><br>Analogous inequalities hold also for $G_{29}$ , $G_{30}, T_{28}, T_{29}, T_{30}$ | 232 |
| It is now sufficient to take $\frac{(a_i)^{(5)}}{(\widehat{M}_{28})^{(5)}}$ , $\frac{(b_i)^{(5)}}{(\widehat{M}_{28})^{(5)}} < 1$ and to choose<br><br>$(\widehat{P}_{28})^{(5)}$ and $(\widehat{Q}_{28})^{(5)}$ large to have | 233 |

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| $\frac{(a_i)^{(5)}}{(\widehat{M}_{28})^{(5)}}\left[(\widehat{P}_{28})^{(5)}+\left((\hat{P}_{28})^{(5)}+G_j^0\right)e^{-\left(\frac{(\hat{P}_{28})^{(5)}+G_j^0}{G_j^0}\right)}\right]\leq(\hat{P}_{28})^{(5)}$ | 234 |
| $\frac{(b_i)^{(5)}}{(\widehat{M}_{28})^{(5)}}\left[\left((\hat{Q}_{28})^{(5)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{28})^{(5)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{28})^{(5)}\right]\leq(\hat{Q}_{28})^{(5)}$ | 235 |
| In order that the operator $\mathcal{A}^{(5)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(5)}$ is a contraction with respect to the metric<br>$d\left(((G_{31})^{(1)},(T_{31})^{(1)}),((G_{31})^{(2)},(T_{31})^{(2)})\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{28})^{(5)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{28})^{(5)}t}\}$<br>Indeed if we denote<br>**Definition of** $(\widetilde{G_{31}}),(\widetilde{T_{31}})$ : $\left((\widetilde{G_{31}}),(\widetilde{T_{31}})\right)=\mathcal{A}^{(5)}\left((G_{31}),(T_{31})\right)$<br>It results<br>$\left\|\tilde{G}_{28}^{(1)}-\tilde{G}_i^{(2)}\right\|\leq\int_0^t(a_{28})^{(5)}\left\|G_{29}^{(1)}-G_{29}^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{(\widehat{M}_{28})^{(5)}s_{(28)}}ds_{(28)}+$<br>$\int_0^t\{(a_{28}')^{(5)}\left\|G_{28}^{(1)}-G_{28}^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}+$<br>$(a_{28}'')^{(5)}\left(T_{29}^{(1)},s_{(28)}\right)\left\|G_{28}^{(1)}-G_{28}^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{(\widehat{M}_{28})^{(5)}s_{(28)}}+$<br>$G_{28}^{(2)}\left\|(a_{28}'')^{(5)}\left(T_{29}^{(1)},s_{(28)}\right)-(a_{28}'')^{(5)}\left(T_{29}^{(2)},s_{(28)}\right)\right\|\ e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{(\widehat{M}_{28})^{(5)}s_{(28)}}\}ds_{(28)}$<br>Where $s_{(28)}$ represents integrand that is integrated over the interval $[0,t]$<br>From the hypotheses on it follows | 236 |
| $\left\|(G_{31})^{(1)}-(G_{31})^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}t}\leq$<br>$\frac{1}{(\widehat{M}_{28})^{(5)}}\left((a_{28})^{(5)}+(a_{28}')^{(5)}+(\widehat{A}_{28})^{(5)}+(\widehat{P}_{28})^{(5)}(\widehat{k}_{28})^{(5)}\right)d\left(((G_{31})^{(1)},(T_{31})^{(1)};\ (G_{31})^{(2)},(T_{31})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 237 |
| **Remark 21:** The fact that we supposed $(a_{28}'')^{(5)}$ and $(b_{28}'')^{(5)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{28})^{(5)}e^{(\widehat{M}_{28})^{(5)}t}$ *and* $(\widehat{Q}_{28})^{(5)}e^{(\widehat{M}_{28})^{(5)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it | 238 |

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| suffices to consider that $(a_i^{''})^{(5)}$ and $(b_i^{''})^{(5)}, i = 28,29,30$ depend only on $T_{29}$ and respectively on $(G_{31})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 22:** There does not exist any $t$ where $G_i\,(t) = 0\ and\ T_i\,(t) = 0$<br><br>it results<br><br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(5)} - (a_i^{''})^{(5)}(T_{29}(s_{(28)}), s_{(28)})\} ds_{(28)}\right]} \geq 0$<br><br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(5)} t\right)} > 0$ for t > 0 | 239 |
| **Definition of** $\left((\widehat{M}_{28})^{(5)}\right)_1, \left((\widehat{M}_{28})^{(5)}\right)_2\ and\ \left((\widehat{M}_{28})^{(5)}\right)_3$ :<br><br>**Remark 23:** if $G_{28}$ is bounded, the same property have also $G_{29}\ and\ G_{30}$ . indeed if<br><br>$G_{28} < (\widehat{M}_{28})^{(5)}$ it follows $\frac{dG_{29}}{dt} \leq \left((\widehat{M}_{28})^{(5)}\right)_1 - (a_{29}^{'})^{(5)} G_{29}$ and by integrating<br><br>$G_{29} \leq \left((\widehat{M}_{28})^{(5)}\right)_2 = G_{29}^0 + 2(a_{29})^{(5)}\left((\widehat{M}_{28})^{(5)}\right)_1 / (a_{29}^{'})^{(5)}$<br><br>In the same way , one can obtain<br><br>$G_{30} \leq \left((\widehat{M}_{28})^{(5)}\right)_3 = G_{30}^0 + 2(a_{30})^{(5)}\left((\widehat{M}_{28})^{(5)}\right)_2 / (a_{30}^{'})^{(5)}$<br><br>If $G_{29}\ or\ G_{30}$ is bounded, the same property follows for $G_{28}$ , $G_{30}$ and $G_{28}$ , $G_{29}$ respectively. | 240 |
| **Remark 24:** If $G_{28}$ *is* bounded, from below, the same property holds for $G_{29}\ and\ G_{30}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{29}$ is bounded from below. | 241 |
| **Remark 25:** If $T_{28}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(5)}((G_{31})(t), t)) = (b_{29}^{'})^{(5)}$ then $T_{29} \to \infty$.<br><br>**Definition of** $(m)^{(5)}$ and $\varepsilon_5$ :<br><br>Indeed let $t_5$ be so that for $t > t_5$<br><br>$(b_{29})^{(5)} - (b_i^{''})^{(5)}((G_{31})(t), t) < \varepsilon_5, T_{28}\,(t) > (m)^{(5)}$ | 242 |
| Then $\frac{dT_{29}}{dt} \geq (a_{29})^{(5)}(m)^{(5)} - \varepsilon_5 T_{29}$ which leads to<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{\varepsilon_5}\right)(1 - e^{-\varepsilon_5 t}) + T_{29}^0 e^{-\varepsilon_5 t}$ If we take t such that $e^{-\varepsilon_5 t} = \frac{1}{2}$ it results<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{2}\right), \quad t = log\frac{2}{\varepsilon_5}$ By taking now $\varepsilon_5$ sufficiently small one sees that $T_{29}$ is unbounded.<br>The same property holds for $T_{30}$ if $\lim_{t\to\infty}(b_{30}^{''})^{(5)}\left((G_{31})(t), t\right) = (b_{30}^{'})^{(5)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations<br><br>Analogous inequalities hold also for $G_{33}$ , $G_{34}, T_{32}, T_{33}, T_{34}$ | 243 |

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| It is now sufficient to take $\frac{(a_i)^{(6)}}{(\widehat{M}_{32})^{(6)}}$ , $\frac{(b_i)^{(6)}}{(\widehat{M}_{32})^{(6)}} < 1$ and to choose<br>$(\widehat{P}_{32})^{(6)}$ and $(\widehat{Q}_{32})^{(6)}$ large to have | 244 |
| $\frac{(a_i)^{(6)}}{(\widehat{M}_{32})^{(6)}}\left[(\widehat{P}_{32})^{(6)} + \left((\widehat{P}_{32})^{(6)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{32})^{(6)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{32})^{(6)}$ | 245 |
| $\frac{(b_i)^{(6)}}{(\widehat{M}_{32})^{(6)}}\left[\left((\widehat{Q}_{32})^{(6)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{32})^{(6)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{32})^{(6)}\right] \leq (\widehat{Q}_{32})^{(6)}$ | 246 |
| In order that the operator $\mathcal{A}^{(6)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(6)}$ is a contraction with respect to the metric<br>$d\left(((G_{35})^{(1)}, (T_{35})^{(1)}), ((G_{35})^{(2)}, (T_{35})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{32})^{(6)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{32})^{(6)}t}\}$<br>Indeed if we denote<br>**Definition of** $\widetilde{(G_{35})}, \widetilde{(T_{35})}$ : $\left(\widetilde{(G_{35})}, \widetilde{(T_{35})}\right) = \mathcal{A}^{(6)}((G_{35}), (T_{35}))$<br>It results<br>$\left\|\tilde{G}_{32}^{(1)} - \tilde{G}_i^{(2)}\right\| \leq \int_0^t (a_{32})^{(6)}\left\|G_{33}^{(1)} - G_{33}^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{(\widehat{M}_{32})^{(6)}s_{(32)}}\,ds_{(32)} +$<br>$\int_0^t\{(a'_{32})^{(6)}\left\|G_{32}^{(1)} - G_{32}^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} +$<br>$(a''_{32})^{(6)}\left(T_{33}^{(1)}, s_{(32)}\right)\left\|G_{32}^{(1)} - G_{32}^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{(\widehat{M}_{32})^{(6)}s_{(32)}} +$<br>$G_{32}^{(2)}\left\|(a''_{32})^{(6)}\left(T_{33}^{(1)}, s_{(32)}\right) - (a''_{32})^{(6)}\left(T_{33}^{(2)}, s_{(32)}\right)\right\|\ e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{(\widehat{M}_{32})^{(6)}s_{(32)}}\}ds_{(32)}$<br>Where $s_{(32)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | 247 |
| $\left\|(G_{35})^{(1)} - (G_{35})^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}t} \leq$<br>$\frac{1}{(\widehat{M}_{32})^{(6)}}\left((a_{32})^{(6)} + (a'_{32})^{(6)} + (\widehat{A}_{32})^{(6)} + (\widehat{P}_{32})^{(6)}(\widehat{k}_{32})^{(6)}\right)d\left(((G_{35})^{(1)}, (T_{35})^{(1)}; (G_{35})^{(2)}, (T_{35})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 248 |
| **Remark 26:** The fact that we supposed $(a''_{32})^{(6)}$ and $(b''_{32})^{(6)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate | 249 |

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| condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{32})^{(6)}e^{(\widehat{M}_{32})^{(6)}t}$ *and* $(\widehat{Q}_{32})^{(6)}e^{(\widehat{M}_{32})^{(6)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(6)}$ and $(b_i^{''})^{(6)}, i = 32,33,34$ depend only on $T_{33}$ and respectively on $(G_{35})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 27:** There does not exist any $t$ where $G_i(t) = 0\ and\ T_i(t) = 0$<br><br>it results<br><br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(6)} - (a_i^{''})^{(6)}(T_{33}(s_{(32)}), s_{(32)})\} ds_{(32)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{(-(b_i^{'})^{(6)}t)} > 0$ for t > 0 | 250 |
| **Definition of** $\left((\widehat{M}_{32})^{(6)}\right)_1, \left((\widehat{M}_{32})^{(6)}\right)_2\ and\ \left((\widehat{M}_{32})^{(6)}\right)_3$ :<br><br>**Remark 28:** if $G_{32}$ is bounded, the same property have also $G_{33}\ and\ G_{34}$ . indeed if<br><br>$G_{32} < (\widehat{M}_{32})^{(6)}$ it follows $\frac{dG_{33}}{dt} \leq \left((\widehat{M}_{32})^{(6)}\right)_1 - (a_{33}^{'})^{(6)}G_{33}$ and by integrating<br><br>$G_{33} \leq \left((\widehat{M}_{32})^{(6)}\right)_2 = G_{33}^0 + 2(a_{33})^{(6)}\left((\widehat{M}_{32})^{(6)}\right)_1/(a_{33}^{'})^{(6)}$<br><br>In the same way , one can obtain<br><br>$G_{34} \leq \left((\widehat{M}_{32})^{(6)}\right)_3 = G_{34}^0 + 2(a_{34})^{(6)}\left((\widehat{M}_{32})^{(6)}\right)_2/(a_{34}^{'})^{(6)}$<br><br>If $G_{33}\ or\ G_{34}$ is bounded, the same property follows for $G_{32}$ , $G_{34}$ and $G_{32}$ , $G_{33}$ respectively. | 251 |
| **Remark 29:** If $G_{32}$ *is* bounded, from below, the same property holds for $G_{33}\ and\ G_{34}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{33}$ is bounded from below. | 252 |
| **Remark 30:** If $T_{32}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(6)}((G_{35})(t), t)) = (b_{33}^{'})^{(6)}$ then $T_{33} \to \infty$.<br><br>**Definition of** $(m)^{(6)}$ and $\varepsilon_6$ :<br><br>Indeed let $t_6$ be so that for $t > t_6$<br><br>$(b_{33})^{(6)} - (b_i^{''})^{(6)}((G_{35})(t), t) < \varepsilon_6, T_{32}(t) > (m)^{(6)}$ | 253 |
| Then $\frac{dT_{33}}{dt} \geq (a_{33})^{(6)}(m)^{(6)} - \varepsilon_6 T_{33}$ which leads to<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{\varepsilon_6}\right)(1 - e^{-\varepsilon_6 t}) + T_{33}^0 e^{-\varepsilon_6 t}$ If we take t such that $e^{-\varepsilon_6 t} = \frac{1}{2}$ it results<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{2}\right),\ t = log\frac{2}{\varepsilon_6}$ By taking now $\varepsilon_6$ sufficiently small one sees that $T_{33}$ is unbounded.<br>The same property holds for $T_{34}$ if $\lim_{t\to\infty}(b_{34}^{''})^{(6)}\left((G_{35})(t), t(t), t\right) = (b_{34}^{'})^{(6)}$ | 254 |

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| We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |
| Analogous inequalities hold also for $G_{37}$ , $G_{38}, T_{36}, T_{37}, T_{38}$<br><br>It is now sufficient to take $\frac{(a_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}$ , $\frac{(b_i)^{(7)}}{(\widehat{M}_{36})^{(7)}} < 1$ and to choose<br>$(\widehat{\mathrm{P}}_{36})^{(7)}$ and $(\widehat{\mathrm{Q}}_{36})^{(7)}$ large to have | 255 |
| $\frac{(a_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}\left[(\widehat{P}_{36})^{(7)} + \left((\hat{P}_{36})^{(7)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{36})^{(7)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{36})^{(7)}$ | 256 |
| $\frac{(b_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}\left[\left((\hat{Q}_{36})^{(7)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{36})^{(7)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{36})^{(7)}\right] \leq (\hat{Q}_{36})^{(7)}$ | 257 |
| In order that the operator $\mathcal{A}^{(7)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(7)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{39})^{(1)}, (T_{39})^{(1)}), ((G_{39})^{(2)}, (T_{39})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{36})^{(7)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{36})^{(7)}t}\}$<br><br>Indeed if we denote<br><br>**Definition of** $(\widetilde{G_{39}}), (\widetilde{T_{39}}) : \left((\widetilde{G_{39}}), (\widetilde{T_{39}})\right) = \mathcal{A}^{(7)}((G_{39}), (T_{39}))$<br>It results<br><br>$\left\|\tilde{G}_{36}^{(1)} - \tilde{G}_i^{(2)}\right\| \leq \int_0^t (a_{36})^{(7)}\left\|G_{37}^{(1)} - G_{37}^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{(\widehat{M}_{36})^{(7)}s_{(36)}}\,ds_{(36)} +$<br>$\int_0^t\{(a'_{36})^{(7)}\left\|G_{36}^{(1)} - G_{36}^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{-(\widehat{M}_{36})^{(7)}s_{(36)}} +$<br>$(a''_{36})^{(7)}\left(T_{37}^{(1)}, s_{(36)}\right)\left\|G_{36}^{(1)} - G_{36}^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{(\widehat{M}_{36})^{(7)}s_{(36)}} +$<br>$G_{36}^{(2)}\left\|(a''_{36})^{(7)}\left(T_{37}^{(1)}, s_{(36)}\right) - (a''_{36})^{(7)}\left(T_{37}^{(2)}, s_{(36)}\right)\right\| \; e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{(\widehat{M}_{36})^{(7)}s_{(36)}}\}ds_{(36)}$<br><br>Where $s_{(36)}$ represents integrand that is integrated over the interval $[0, t]$<br><br>From the hypotheses on it follows | 258 |
| $\left\|(G_{39})^{(1)} - (G_{39})^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}t} \leq$<br>$\frac{1}{(\widehat{M}_{36})^{(7)}}\left((a_{36})^{(7)} + (a'_{36})^{(7)} + (\widehat{A}_{36})^{(7)} + (\widehat{P}_{36})^{(7)}(\widehat{k}_{36})^{(7)}\right)d\left(((G_{39})^{(1)}, (T_{39})^{(1)}; (G_{39})^{(2)}, (T_{39})^{(2)})\right)$<br><br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 259 |

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| **Remark 31:** The fact that we supposed $(a_{36}^{''})^{(7)}$ and $(b_{36}^{''})^{(7)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{36})^{(7)}e^{(\widehat{M}_{36})^{(7)}t}$ $and$ $(\widehat{Q}_{36})^{(7)}e^{(\widehat{M}_{36})^{(7)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(7)}$ and $(b_i^{''})^{(7)}, i = 36,37,38$ depend only on $T_{37}$ and respectively on $(G_{39})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 260 |
| **Remark 32:** There does not exist any $t$ where $G_i\,(t) = 0\ and\ T_i\,(t) = 0$<br><br>it results<br><br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t\{(a_i^{'})^{(7)}-(a_i^{''})^{(7)}(T_{37}(s_{(36)}),s_{(36)})\}ds_{(36)}\right]} \geq 0$<br><br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(7)}t\right)} > 0$ for t > 0 | 261 |
| **Definition of** $\left((\widehat{M}_{36})^{(7)}\right)_1, \left((\widehat{M}_{36})^{(7)}\right)_2\ and\ \left((\widehat{M}_{36})^{(7)}\right)_3$ :<br><br>**Remark 33:** if $G_{36}$ is bounded, the same property have also $G_{37}\ and\ G_{38}$ . indeed if<br><br>$G_{36} < (\widehat{M}_{36})^{(7)}$ it follows $\frac{dG_{37}}{dt} \leq \left((\widehat{M}_{36})^{(7)}\right)_1 - (a_{37}^{'})^{(7)}G_{37}$ and by integrating<br><br>$G_{37} \leq \left((\widehat{M}_{36})^{(7)}\right)_2 = G_{37}^0 + 2(a_{37})^{(7)}\left((\widehat{M}_{36})^{(7)}\right)_1/(a_{37}^{'})^{(7)}$<br><br>In the same way , one can obtain<br><br>$G_{38} \leq \left((\widehat{M}_{36})^{(7)}\right)_3 = G_{38}^0 + 2(a_{38})^{(7)}\left((\widehat{M}_{36})^{(7)}\right)_2/(a_{38}^{'})^{(7)}$<br><br>If $G_{37}\ or\ G_{38}$ is bounded, the same property follows for $G_{36}$ , $G_{38}$ and $G_{36}$ , $G_{37}$ respectively. | 262 |
| **Remark 34:** If $G_{36}$ $is$ bounded, from below, the same property holds for $G_{37}\ and\ G_{38}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{37}$ is bounded from below. | 263 |
| **Remark 35:** If $T_{36}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(7)}((G_{39})(t),t)) = (b_{37}^{'})^{(7)}$ then $T_{37} \to \infty$.<br><br>**Definition of** $(m)^{(7)}$ and $\varepsilon_7$ :<br><br>Indeed let $t_7$ be so that for $t > t_7$<br><br>$(b_{37})^{(7)} - (b_i^{''})^{(7)}((G_{39})(t),t) < \varepsilon_7, T_{36}\,(t) > (m)^{(7)}$ | 264 |

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| Then $\frac{dT_{37}}{dt} \geq (a_{37})^{(7)}(m)^{(7)} - \varepsilon_7 T_{37}$ which leads to<br><br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{\varepsilon_7}\right)(1 - e^{-\varepsilon_7 t}) + T_{37}^0 e^{-\varepsilon_7 t}$ If we take t such that $e^{-\varepsilon_7 t} = \frac{1}{2}$ it results<br><br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{2}\right), \quad t = log\frac{2}{\varepsilon_7}$ By taking now $\varepsilon_7$ sufficiently small one sees that $T_{37}$ is unbounded.<br>The same property holds for $T_{38}$ if $\lim_{t\to\infty}(b_{38}^{''})^{(7)}\left((G_{39})(t), t\right) = (b_{38}^{'})^{(7)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 265 |
| It is now sufficient to take $\frac{(a_i)^{(8)}}{(\bar{M}_{40})^{(8)}}$ , $\frac{(b_i)^{(8)}}{(\bar{M}_{40})^{(8)}} < 1$ and to choose<br>$(\hat{P}_{40})^{(8)}$ and $(\hat{Q}_{40})^{(8)}$ large to have | 266 |
| $\frac{(a_i)^{(8)}}{(\bar{M}_{40})^{(8)}}\left[(\hat{P}_{40})^{(8)} + \left((\hat{P}_{40})^{(8)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{40})^{(8)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{40})^{(8)}$ | 267 |
| $\frac{(b_i)^{(8)}}{(\bar{M}_{40})^{(8)}}\left[\left((\hat{Q}_{40})^{(8)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{40})^{(8)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{40})^{(8)}\right] \leq (\hat{Q}_{40})^{(8)}$ | 268 |
| In order that the operator $\mathcal{A}^{(8)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(8)}$ is a contraction with respect to the metric | |
| $d\left(((G_{43})^{(1)}, (T_{43})^{(1)}), ((G_{43})^{(2)}, (T_{43})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\bar{M}_{40})^{(8)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\bar{M}_{40})^{(8)}t}\}$ | 269 |
| Indeed if we denote<br><br>**Definition of** $\widetilde{(G_{43})}, \widetilde{(T_{43})}$ : $\left(\widetilde{(G_{43})}, \widetilde{(T_{43})}\right) = \mathcal{A}^{(8)}((G_{43}), (T_{43}))$ | 270 |
| It results<br><br>$\left\|\tilde{G}_{40}^{(1)} - \tilde{G}_i^{(2)}\right\| \leq \int_0^t (a_{40})^{(8)}\left\|G_{41}^{(1)} - G_{41}^{(2)}\right\|e^{-(\bar{M}_{40})^{(8)}s_{(40)}}e^{(\bar{M}_{40})^{(8)}s_{(40)}}\,ds_{(40)} +$<br>$\int_0^t\{(a_{40}^{'})^{(8)}\left\|G_{40}^{(1)} - G_{40}^{(2)}\right\|e^{-(\bar{M}_{40})^{(8)}s_{(40)}}e^{-(\bar{M}_{40})^{(8)}s_{(40)}} +$<br>$(a_{40}^{''})^{(8)}\left(T_{41}^{(1)}, s_{(40)}\right)\left\|G_{40}^{(1)} - G_{40}^{(2)}\right\|e^{-(\bar{M}_{40})^{(8)}s_{(40)}}e^{(\bar{M}_{40})^{(8)}s_{(40)}} +$<br>$G_{40}^{(2)}\left\|(a_{40}^{''})^{(8)}\left(T_{41}^{(1)}, s_{(40)}\right) - (a_{40}^{''})^{(8)}\left(T_{41}^{(2)}, s_{(40)}\right)\right\| e^{-(\bar{M}_{40})^{(8)}s_{(40)}}e^{(\bar{M}_{40})^{(8)}s_{(40)}}\}ds_{(40)}$ | 271 |
| Where $s_{(40)}$ represents integrand that is integrated over the interval [0, t] | 272 |

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| From the hypotheses it follows | |
| $\left\|(G_{43})^{(1)} - (G_{43})^{(2)}\right\| e^{-(\widehat{M}_{40})^{(8)}t} \leq$ $\frac{1}{(\widehat{M}_{40})^{(8)}}\left((a_{40})^{(8)} + (a'_{40})^{(8)} + (\widehat{A}_{40})^{(8)} + (\widehat{P}_{40})^{(8)}(\widehat{k}_{40})^{(8)}\right) d\left(((G_{43})^{(1)}, (T_{43})^{(1)};\ (G_{43})^{(2)}, (T_{43})^{(2)})\right)$ And analogous inequalities for $G_i$ $and$ $T_i$. Taking into account the hypothesis the result follows | 273 |
| **Remark 36:** The fact that we supposed $(a''_{40})^{(8)}$ and $(b''_{40})^{(8)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{40})^{(8)} e^{(\widehat{M}_{40})^{(8)}t}$ $and$ $(\widehat{Q}_{40})^{(8)} e^{(\widehat{M}_{40})^{(8)}t}$ respectively of $\mathbb{R}_+$. If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a''_i)^{(8)}$ and $(b''_i)^{(8)}, i = 40,41,42$ depend only on $T_{41}$ and respectively on $(G_{43})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 274 |
| **Remark 37** There does not exist any $t$ where $G_i(t) = 0$ $and$ $T_i(t) = 0$ it results $G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a'_i)^{(8)} - (a''_i)^{(8)}(T_{41}(s_{(40)}), s_{(40)})\} ds_{(40)}\right]} \geq 0$ $T_i(t) \geq T_i^0 e^{\left(-(b'_i)^{(8)}t\right)} > 0$ for t > 0 | 275 |
| **Definition of** $\left((\widehat{M}_{40})^{(8)}\right)_1, \left((\widehat{M}_{40})^{(8)}\right)_2$ $and$ $\left((\widehat{M}_{40})^{(8)}\right)_3$ : **Remark 38:** if $G_{40}$ is bounded, the same property have also $G_{41}$ $and$ $G_{42}$ . indeed if $G_{40} < (\widehat{M}_{40})^{(8)}$ it follows $\frac{dG_{41}}{dt} \leq \left((\widehat{M}_{40})^{(8)}\right)_1 - (a'_{41})^{(8)} G_{41}$ and by integrating $G_{41} \leq \left((\widehat{M}_{40})^{(8)}\right)_2 = G_{41}^0 + 2(a_{41})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_1 / (a'_{41})^{(8)}$ In the same way , one can obtain $G_{42} \leq \left((\widehat{M}_{40})^{(8)}\right)_3 = G_{42}^0 + 2(a_{42})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_2 / (a'_{42})^{(8)}$ If $G_{41}$ $or$ $G_{42}$ is bounded, the same property follows for $G_{40}$ , $G_{42}$ and $G_{40}$ , $G_{41}$ respectively. | 276 |
| **Remark 39:** If $G_{40}$ $is$ bounded, from below, the same property holds for $G_{41}$ $and$ $G_{42}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{41}$ is bounded from below. | 277 |
| **Remark 40:** If $T_{40}$ is bounded from below and $\lim_{t\to\infty}((b''_i)^{(8)}((G_{43})(t), t)) = (b'_{41})^{(8)}$ then $T_{41} \to \infty$. | 278 |

**Definition of** $(m)^{(8)}$ and $\varepsilon_8$ :

Indeed let $t_8$ be so that for $t > t_8$

$(b_{41})^{(8)} - (b_i^{''})^{(8)}\big((G_{43})(t), t\big) < \varepsilon_8, T_{40}(t) > (m)^{(8)}$

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Then $\frac{dT_{41}}{dt} \geq (a_{41})^{(8)}(m)^{(8)} - \varepsilon_8 T_{41}$ which leads to

$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{\varepsilon_8}\right)(1 - e^{-\varepsilon_8 t}) + T_{41}^0 e^{-\varepsilon_8 t}$ If we take t such that $e^{-\varepsilon_8 t} = \frac{1}{2}$ it results

$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{2}\right), \quad t = log\frac{2}{\varepsilon_8}$ By taking now $\varepsilon_8$ sufficiently small one sees that $T_{41}$ is unbounded.
The same property holds for $T_{42}$ if $\lim_{t\to\infty}(b_{42}^{''})^{(8)}\big((G_{43})(t), t(t), t\big) = (b_{42}^{'})^{(8)}$

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It is now sufficient to take $\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}, \frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}} < 1$ and to choose $(\widehat{P}_{44})^{(9)}$ and $(\widehat{Q}_{44})^{(9)}$ large to have

$$\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[(\widehat{P}_{44})^{(9)} + \big((\hat{P}_{44})^{(9)} + G_j^0\big)e^{-\left(\frac{(\hat{P}_{44})^{(9)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{44})^{(9)}$$

$$\frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[\big((\hat{Q}_{44})^{(9)} + T_j^0\big)e^{-\left(\frac{(\hat{Q}_{44})^{(9)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{44})^{(9)}\right] \leq (\hat{Q}_{44})^{(9)}$$

In order that the operator $\mathcal{A}^{(9)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying 39,35,36 into itself

The operator $\mathcal{A}^{(9)}$ is a contraction with respect to the metric

$$d\Big(\big((G_{47})^{(1)}, (T_{47})^{(1)}\big), \big((G_{47})^{(2)}, (T_{47})^{(2)}\big)\Big) =$$
$$sup_i\{max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t) - G_i^{(2)}(t)\right|e^{-(\widehat{M}_{44})^{(9)}t}, max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t) - T_i^{(2)}(t)\right|e^{-(\widehat{M}_{44})^{(9)}t}\}$$

Indeed if we denote

**Definition of** $\widetilde{(G_{47})}, \widetilde{(T_{47})} : \big(\widetilde{(G_{47})}, \widetilde{(T_{47})}\big) = \mathcal{A}^{(9)}\big((G_{47}), (T_{47})\big)$
It results

$$\left|\tilde{G}_{44}^{(1)} - \tilde{G}_i^{(2)}\right| \leq \int_0^t (a_{44})^{(9)}\left|G_{45}^{(1)} - G_{45}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\,ds_{(44)} +$$
$$\int_0^t\{(a_{44}^{'})^{(9)}\left|G_{44}^{(1)} - G_{44}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{-(\widehat{M}_{44})^{(9)}s_{(44)}} +$$
$$(a_{44}^{''})^{(9)}\big(T_{45}^{(1)}, s_{(44)}\big)\left|G_{44}^{(1)} - G_{44}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}} +$$
$$G_{44}^{(2)}\left|(a_{44}^{''})^{(9)}\big(T_{45}^{(1)}, s_{(44)}\big) - (a_{44}^{''})^{(9)}\big(T_{45}^{(2)}, s_{(44)}\big)\right|\ e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\}ds_{(44)}$$

Where $s_{(44)}$ represents integrand that is integrated over the interval $[0, t]$

From the hypotheses on 45,46,47,28 and 29 it follows

$$\left|(G_{47})^{(1)} - G^{(2)}\right| e^{-(\widehat{M}_{44})^{(9)}t} \leq \frac{1}{(\widehat{M}_{44})^{(9)}}\left((a_{44})^{(9)} + (a'_{44})^{(9)} + (\widehat{A}_{44})^{(9)} + (\widehat{P}_{44})^{(9)}(\widehat{k}_{44})^{(9)}\right) d\left(((G_{47})^{(1)}, (T_{47})^{(1)};\ (G_{47})^{(2)}, (T_{47})^{(2)})\right)$$

And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis (39,35,36) the result follows

**Remark 41:** The fact that we supposed $(a''_{44})^{(9)}$ and $(b''_{44})^{(9)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{44})^{(9)} e^{(\widehat{M}_{44})^{(9)}t}$ *and* $(\widehat{Q}_{44})^{(9)} e^{(\widehat{M}_{44})^{(9)}t}$ respectively of $\mathbb{R}_+$.

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i'')^{(9)}$ and $(b_i'')^{(9)}, i = 44,45,46$ depend only on $T_{45}$ and respectively on $(G_{47})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition.

**Remark 42:** There does not exist any $t$ where $G_i(t) = 0\ and\ T_i(t) = 0$

From 99 to 44 it results

$$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a'_i)^{(9)} - (a''_i)^{(9)}(T_{45}(s_{(44)}), s_{(44)})\} ds_{(44)}\right]} \geq 0$$

$$T_i(t) \geq T_i^0 e^{(-(b'_i)^{(9)}t)} > 0 \quad \text{for t} > 0$$

**Definition of** $\left((\widehat{M}_{44})^{(9)}\right)_1, \left((\widehat{M}_{44})^{(9)}\right)_2$ *and* $\left((\widehat{M}_{44})^{(9)}\right)_3$ :

**Remark 43:** if $G_{44}$ is bounded, the same property have also $G_{45}\ and\ G_{46}$ . indeed if

$G_{44} < (\widehat{M}_{44})^{(9)}$ it follows $\frac{dG_{45}}{dt} \leq \left((\widehat{M}_{44})^{(9)}\right)_1 - (a'_{45})^{(9)} G_{45}$ and by integrating

$$G_{45} \leq \left((\widehat{M}_{44})^{(9)}\right)_2 = G_{45}^0 + 2(a_{45})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_1 / (a'_{45})^{(9)}$$

In the same way , one can obtain

$$G_{46} \leq \left((\widehat{M}_{44})^{(9)}\right)_3 = G_{46}^0 + 2(a_{46})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_2 / (a'_{46})^{(9)}$$

If $G_{45}\ or\ G_{46}$ is bounded, the same property follows for $G_{44}$ , $G_{46}$ and $G_{44}$ , $G_{45}$ respectively.

**Remark 44:** If $G_{44}$ *is* bounded, from below, the same property holds for $G_{45}\ and\ G_{46}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{45}$ is bounded from below.

**Remark 45:** If $T_{44}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(9)}\left((G_{47})(t), t\right)) = (b'_{45})^{(9)}$ then $T_{45} \to \infty$.

**Definition of** $(m)^{(9)}$ and $\varepsilon_9$ :

Indeed let $t_9$ be so that for $t > t_9$

$$(b_{45})^{(9)} - (b_i^{''})^{(9)}\big((G_{47})(t),t\big) < \varepsilon_9, T_{44}\,(t) > (m)^{(9)}$$

Then $\frac{dT_{45}}{dt} \geq (a_{45})^{(9)}(m)^{(9)} - \varepsilon_9 T_{45}$ which leads to

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{\varepsilon_9}\right)(1 - e^{-\varepsilon_9 t}) + T_{45}^0 e^{-\varepsilon_9 t}$ If we take t such that $e^{-\varepsilon_9 t} = \frac{1}{2}$ it results

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{2}\right), \quad t = log\frac{2}{\varepsilon_9}$ By taking now $\varepsilon_9$ sufficiently small one sees that $T_{45}$ is unbounded.
The same property holds for $T_{46}$ if $\lim_{t\to\infty}(b_{46}^{''})^{(9)}\left((G_{47})(t),t\right) = (b_{46}^{'})^{(9)}$
We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 92

**Behavior of the solutions of equation** 280

**Theorem** If we denote and define

**Definition of** $(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ :

$(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ four constants satisfying
$-(\sigma_2)^{(1)} \leq -(a_{13}^{'})^{(1)} + (a_{14}^{'})^{(1)} - (a_{13}^{''})^{(1)}(T_{14}, t) + (a_{14}^{''})^{(1)}(T_{14}, t) \leq -(\sigma_1)^{(1)}$

$-(\tau_2)^{(1)} \leq -(b_{13}^{'})^{(1)} + (b_{14}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G,t) - (b_{14}^{''})^{(1)}(G,t) \leq -(\tau_1)^{(1)}$

**Definition of** $(\nu_1)^{(1)}, (\nu_2)^{(1)}, (u_1)^{(1)}, (u_2)^{(1)}, \nu^{(1)}, u^{(1)}$ : 281

By $(\nu_1)^{(1)} > 0$, $(\nu_2)^{(1)} < 0$ and respectively $(u_1)^{(1)} > 0$, $(u_2)^{(1)} < 0$ the roots of the equations
$(a_{14})^{(1)}\left(\nu^{(1)}\right)^2 + (\sigma_1)^{(1)}\nu^{(1)} - (a_{13})^{(1)} = 0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2 + (\tau_1)^{(1)}u^{(1)} - (b_{13})^{(1)} = 0$

**Definition of** $(\bar{\nu}_1)^{(1)}, , (\bar{\nu}_2)^{(1)}, (\bar{u}_1)^{(1)}, (\bar{u}_2)^{(1)}$ : 282

By $(\bar{\nu}_1)^{(1)} > 0$, $(\bar{\nu}_2)^{(1)} < 0$ and respectively $(\bar{u}_1)^{(1)} > 0$, $(\bar{u}_2)^{(1)} < 0$ the roots of the equations
$(a_{14})^{(1)}\left(\nu^{(1)}\right)^2 + (\sigma_2)^{(1)}\nu^{(1)} - (a_{13})^{(1)} = 0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2 + (\tau_2)^{(1)}u^{(1)} - (b_{13})^{(1)} = 0$

**Definition of** $(m_1)^{(1)}, (m_2)^{(1)}, (\mu_1)^{(1)}, (\mu_2)^{(1)}, (\nu_0)^{(1)}$ :- 283

If we define $(m_1)^{(1)}, (m_2)^{(1)}, (\mu_1)^{(1)}, (\mu_2)^{(1)}$ by
$(m_2)^{(1)} = (\nu_0)^{(1)}, (m_1)^{(1)} = (\nu_1)^{(1)}, \; if\; (\nu_0)^{(1)} < (\nu_1)^{(1)}$

$(m_2)^{(1)} = (\nu_1)^{(1)}, (m_1)^{(1)} = (\bar{\nu}_1)^{(1)}, if\; (\nu_1)^{(1)} < (\nu_0)^{(1)} < (\bar{\nu}_1)^{(1)},$

and $\boxed{(\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}}$

$(m_2)^{(1)} = (\nu_1)^{(1)}, (m_1)^{(1)} = (\nu_0)^{(1)}, \; if\; (\bar{\nu}_1)^{(1)} < (\nu_0)^{(1)}$

and analogously 284

$(\mu_2)^{(1)} = (u_0)^{(1)}, (\mu_1)^{(1)} = (u_1)^{(1)}, \; if\; (u_0)^{(1)} < (u_1)^{(1)}$

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| $(\mu_2)^{(1)} = (u_1)^{(1)}, (\mu_1)^{(1)} = (\bar{u}_1)^{(1)}, if\ (u_1)^{(1)} < (u_0)^{(1)} < (\bar{u}_1)^{(1)},$<br>and $\boxed{(u_0)^{(1)} = \frac{T_{13}^0}{T_{14}^0}}$<br>$(\mu_2)^{(1)} = (u_1)^{(1)}, (\mu_1)^{(1)} = (u_0)^{(1)}, if\ (\bar{u}_1)^{(1)} < (u_0)^{(1)}$ where $(u_1)^{(1)}, (\bar{u}_1)^{(1)}$<br>are defined | |
| Then the solution of global equations satisfies the inequalities<br>$G_{13}^0 e^{((S_1)^{(1)} - (p_{13})^{(1)})t} \le G_{13}(t) \le G_{13}^0 e^{(S_1)^{(1)} t}$<br>where $(p_i)^{(1)}$ is defined by equation<br>$\frac{1}{(m_1)^{(1)}} G_{13}^0 e^{((S_1)^{(1)} - (p_{13})^{(1)})t} \le G_{14}(t) \le \frac{1}{(m_2)^{(1)}} G_{13}^0 e^{(S_1)^{(1)} t}$ | 285 |
| $\left(\frac{(a_{15})^{(1)} G_{13}^0}{(m_1)^{(1)}((S_1)^{(1)} - (p_{13})^{(1)} - (S_2)^{(1)})}\left[e^{((S_1)^{(1)} - (p_{13})^{(1)})t} - e^{-(S_2)^{(1)} t}\right] + G_{15}^0 e^{-(S_2)^{(1)} t} \le G_{15}(t) \le \right.$<br>$\left.\frac{(a_{15})^{(1)} G_{13}^0}{(m_2)^{(1)}((S_1)^{(1)} - (a'_{15})^{(1)})}[e^{(S_1)^{(1)} t} - e^{-(a'_{15})^{(1)} t}] + G_{15}^0 e^{-(a'_{15})^{(1)} t}\right)$ | 286 |
| $\boxed{T_{13}^0 e^{(R_1)^{(1)} t} \le T_{13}(t) \le T_{13}^0 e^{((R_1)^{(1)} + (r_{13})^{(1)})t}}$ | 287 |
| $\frac{1}{(\mu_1)^{(1)}} T_{13}^0 e^{(R_1)^{(1)} t} \le T_{13}(t) \le \frac{1}{(\mu_2)^{(1)}} T_{13}^0 e^{((R_1)^{(1)} + (r_{13})^{(1)})t}$ | 288 |
| $\frac{(b_{15})^{(1)} T_{13}^0}{(\mu_1)^{(1)}((R_1)^{(1)} - (b'_{15})^{(1)})}\left[e^{(R_1)^{(1)} t} - e^{-(b'_{15})^{(1)} t}\right] + T_{15}^0 e^{-(b'_{15})^{(1)} t} \le T_{15}(t) \le$<br>$\frac{(a_{15})^{(1)} T_{13}^0}{(\mu_2)^{(1)}((R_1)^{(1)} + (r_{13})^{(1)} + (R_2)^{(1)})}\left[e^{((R_1)^{(1)} + (r_{13})^{(1)})t} - e^{-(R_2)^{(1)} t}\right] + T_{15}^0 e^{-(R_2)^{(1)} t}$ | 289 |
| **Definition of** $(S_1)^{(1)}, (S_2)^{(1)}, (R_1)^{(1)}, (R_2)^{(1)}$:-<br>Where $(S_1)^{(1)} = (a_{13})^{(1)}(m_2)^{(1)} - (a'_{13})^{(1)}$<br>$(S_2)^{(1)} = (a_{15})^{(1)} - (p_{15})^{(1)}$<br>$(R_1)^{(1)} = (b_{13})^{(1)}(\mu_2)^{(1)} - (b'_{13})^{(1)}$<br>$(R_2)^{(1)} = (b'_{15})^{(1)} - (r_{15})^{(1)}$ | 290 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define | 291 |
| **Definition of** $(\sigma_1)^{(2)}, (\sigma_2)^{(2)}, (\tau_1)^{(2)}, (\tau_2)^{(2)}$ :<br>$(\sigma_1)^{(2)}, (\sigma_2)^{(2)}, (\tau_1)^{(2)}, (\tau_2)^{(2)}$ four constants satisfying | 292 |

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| $-(\sigma_2)^{(2)} \le -(a_{16}^{'})^{(2)} + (a_{17}^{'})^{(2)} - (a_{16}^{''})^{(2)}(\mathrm{T}_{17}, t) + (a_{17}^{''})^{(2)}(\mathrm{T}_{17}, t) \le -(\sigma_1)^{(2)}$ | 293 |
| $-(\tau_2)^{(2)} \le -(b_{16}^{'})^{(2)} + (b_{17}^{'})^{(2)} - (b_{16}^{''})^{(2)}\big((G_{19}), t\big) - (b_{17}^{''})^{(2)}\big((G_{19}), t\big) \le -(\tau_1)^{(2)}$ | 294 |
| **Definition of** $(\nu_1)^{(2)}, (\nu_2)^{(2)}, (u_1)^{(2)}, (u_2)^{(2)}$ : | 295 |
| By $(\nu_1)^{(2)} > 0$ , $(\nu_2)^{(2)} < 0$ and respectively $(u_1)^{(2)} > 0$ , $(u_2)^{(2)} < 0$ the roots | 296 |
| of the equations $(a_{17})^{(2)}\big(\nu^{(2)}\big)^2 + (\sigma_1)^{(2)}\nu^{(2)} - (a_{16})^{(2)} = 0$ | 297 |
| and $(b_{14})^{(2)}\big(u^{(2)}\big)^2 + (\tau_1)^{(2)}u^{(2)} - (b_{16})^{(2)} = 0$ and | 298 |
| **Definition of** $(\bar{\nu}_1)^{(2)}, , (\bar{\nu}_2)^{(2)}, (\bar{u}_1)^{(2)}, (\bar{u}_2)^{(2)}$ : | 299 |
| By $(\bar{\nu}_1)^{(2)} > 0$ , $(\bar{\nu}_2)^{(2)} < 0$ and respectively $(\bar{u}_1)^{(2)} > 0$ , $(\bar{u}_2)^{(2)} < 0$ the | 300 |
| roots of the equations $(a_{17})^{(2)}\big(\nu^{(2)}\big)^2 + (\sigma_2)^{(2)}\nu^{(2)} - (a_{16})^{(2)} = 0$ | 301 |
| and $(b_{17})^{(2)}\big(u^{(2)}\big)^2 + (\tau_2)^{(2)}u^{(2)} - (b_{16})^{(2)} = 0$ | 302 |
| **Definition of** $(m_1)^{(2)}$ , $(m_2)^{(2)}$ , $(\mu_1)^{(2)}, (\mu_2)^{(2)}$ :- | 303 |
| If we define $(m_1)^{(2)}$ , $(m_2)^{(2)}$ , $(\mu_1)^{(2)}, (\mu_2)^{(2)}$ by | 304 |
| $(m_2)^{(2)} = (\nu_0)^{(2)}, (m_1)^{(2)} = (\nu_1)^{(2)}$, ***if*** $(\nu_0)^{(2)} < (\nu_1)^{(2)}$ | 305 |
| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ , ***if*** $(\nu_1)^{(2)} < (\nu_0)^{(2)} < (\bar{\nu}_1)^{(2)}$, and $\boxed{(\nu_0)^{(2)} = \frac{\mathrm{G}_{16}^0}{\mathrm{G}_{17}^0}}$ | 306 |
| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\nu_0)^{(2)}$, ***if*** $(\bar{\nu}_1)^{(2)} < (\nu_0)^{(2)}$ | 307 |
| and analogously<br>$(\mu_2)^{(2)} = (u_0)^{(2)}, (\mu_1)^{(2)} = (u_1)^{(2)}$, ***if*** $(u_0)^{(2)} < (u_1)^{(2)}$<br>$(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (\bar{u}_1)^{(2)}$ , ***if*** $(u_1)^{(2)} < (u_0)^{(2)} < (\bar{u}_1)^{(2)}$, and $\boxed{(u_0)^{(2)} = \frac{\mathrm{T}_{16}^0}{\mathrm{T}_{17}^0}}$ | 308 |
| $(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (u_0)^{(2)}$, ***if*** $(\bar{u}_1)^{(2)} < (u_0)^{(2)}$ | 309 |
| Then the solution of global equations satisfies the inequalities<br>$\mathrm{G}_{16}^0 \mathrm{e}^{((\mathrm{S}_1)^{(2)} - (p_{16})^{(2)})\mathrm{t}} \le G_{16}(t) \le \mathrm{G}_{16}^0 \mathrm{e}^{(\mathrm{S}_1)^{(2)}\mathrm{t}}$ | 310 |
| $(p_i)^{(2)}$ is defined by equation | |

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| $\frac{1}{(m_1)^{(2)}} \mathrm{G}_{16}^{0} \mathrm{e}^{((\mathrm{S}_1)^{(2)}-(p_{16})^{(2)})\mathrm{t}} \leq G_{17}(t) \leq \frac{1}{(m_2)^{(2)}} \mathrm{G}_{16}^{0} \mathrm{e}^{(\mathrm{S}_1)^{(2)}\mathrm{t}}$ | 311 |
| $\left(\frac{(a_{18})^{(2)}\mathrm{G}_{16}^{0}}{(m_1)^{(2)}((\mathrm{S}_1)^{(2)}-(p_{16})^{(2)}-(\mathrm{S}_2)^{(2)})}\left[\mathrm{e}^{((\mathrm{S}_1)^{(2)}-(p_{16})^{(2)})\mathrm{t}} - \mathrm{e}^{-(\mathrm{S}_2)^{(2)}\mathrm{t}}\right] + \mathrm{G}_{18}^{0}\mathrm{e}^{-(\mathrm{S}_2)^{(2)}\mathrm{t}} \leq \mathrm{G}_{18}(t) \leq \right.$ $\left.\frac{(a_{18})^{(2)}\mathrm{G}_{16}^{0}}{(m_2)^{(2)}((\mathrm{S}_1)^{(2)}-(a_{18}^{'})^{(2)})}\left[\mathrm{e}^{(\mathrm{S}_1)^{(2)}\mathrm{t}} - \mathrm{e}^{-(a_{18}^{'})^{(2)}\mathrm{t}}\right] + \mathrm{G}_{18}^{0}\mathrm{e}^{-(a_{18}^{'})^{(2)}\mathrm{t}}\right)$ | 312 |
| $\boxed{\mathrm{T}_{16}^{0}\mathrm{e}^{(\mathrm{R}_1)^{(2)}t} \leq T_{16}(t) \leq \mathrm{T}_{16}^{0}\mathrm{e}^{((\mathrm{R}_1)^{(2)}+(r_{16})^{(2)})t}}$ | 313 |
| $\frac{1}{(\mu_1)^{(2)}} \mathrm{T}_{16}^{0}\mathrm{e}^{(\mathrm{R}_1)^{(2)}t} \leq T_{16}(t) \leq \frac{1}{(\mu_2)^{(2)}} \mathrm{T}_{16}^{0}\mathrm{e}^{((\mathrm{R}_1)^{(2)}+(r_{16})^{(2)})t}$ | 314 |
| $\frac{(b_{18})^{(2)}\mathrm{T}_{16}^{0}}{(\mu_1)^{(2)}((\mathrm{R}_1)^{(2)}-(b_{18}^{'})^{(2)})}\left[\mathrm{e}^{(\mathrm{R}_1)^{(2)}t} - \mathrm{e}^{-(b_{18}^{'})^{(2)}t}\right] + \mathrm{T}_{18}^{0}\mathrm{e}^{-(b_{18}^{'})^{(2)}t} \leq T_{18}(t) \leq$ $\frac{(a_{18})^{(2)}\mathrm{T}_{16}^{0}}{(\mu_2)^{(2)}((\mathrm{R}_1)^{(2)}+(r_{16})^{(2)}+(\mathrm{R}_2)^{(2)})}\left[\mathrm{e}^{((\mathrm{R}_1)^{(2)}+(r_{16})^{(2)})t} - \mathrm{e}^{-(\mathrm{R}_2)^{(2)}t}\right] + \mathrm{T}_{18}^{0}\mathrm{e}^{-(\mathrm{R}_2)^{(2)}t}$ | 315 |
| **Definition of** $(\mathrm{S}_1)^{(2)}, (\mathrm{S}_2)^{(2)}, (\mathrm{R}_1)^{(2)}, (\mathrm{R}_2)^{(2)}$:- | 316 |
| Where $(\mathrm{S}_1)^{(2)} = (a_{16})^{(2)}(m_2)^{(2)} - (a_{16}^{'})^{(2)}$ $(\mathrm{S}_2)^{(2)} = (a_{18})^{(2)} - (p_{18})^{(2)}$ | 317 |
| $(R_1)^{(2)} = (b_{16})^{(2)}(\mu_2)^{(1)} - (b_{16}^{'})^{(2)}$ $(\mathrm{R}_2)^{(2)} = (b_{18}^{'})^{(2)} - (r_{18})^{(2)}$ | 318 |
| **Behavior of the solutions** **Theorem 3:** If we denote and define **Definition of** $(\sigma_1)^{(3)}, (\sigma_2)^{(3)}, (\tau_1)^{(3)}, (\tau_2)^{(3)}$: $(\sigma_1)^{(3)}, (\sigma_2)^{(3)}, (\tau_1)^{(3)}, (\tau_2)^{(3)}$ four constants satisfying $-(\sigma_2)^{(3)} \leq -(a_{20}^{'})^{(3)} + (a_{21}^{'})^{(3)} - (a_{20}^{''})^{(3)}(T_{21}, t) + (a_{21}^{''})^{(3)}(T_{21}, t) \leq -(\sigma_1)^{(3)}$ $-(\tau_2)^{(3)} \leq -(b_{20}^{'})^{(3)} + (b_{21}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23}, t) - (b_{21}^{''})^{(3)}((G_{23}), t) \leq -(\tau_1)^{(3)}$ | 319 |
| **Definition of** $(\nu_1)^{(3)}, (\nu_2)^{(3)}, (u_1)^{(3)}, (u_2)^{(3)}$: By $(\nu_1)^{(3)} > 0$, $(\nu_2)^{(3)} < 0$ and respectively $(u_1)^{(3)} > 0$, $(u_2)^{(3)} < 0$ the roots of the equations $(a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_1)^{(3)}\nu^{(3)} - (a_{20})^{(3)} = 0$ and $(b_{21})^{(3)}\left(u^{(3)}\right)^2 + (\tau_1)^{(3)}u^{(3)} - (b_{20})^{(3)} = 0$ and By $(\bar{\nu}_1)^{(3)} > 0$, $(\bar{\nu}_2)^{(3)} < 0$ and respectively $(\bar{u}_1)^{(3)} > 0$, $(\bar{u}_2)^{(3)} < 0$ the | 320 |

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| roots of the equations $(a_{21})^{(3)}\left(v^{(3)}\right)^2 + (\sigma_2)^{(3)}v^{(3)} - (a_{20})^{(3)} = 0$<br>and $(b_{21})^{(3)}\left(u^{(3)}\right)^2 + (\tau_2)^{(3)}u^{(3)} - (b_{20})^{(3)} = 0$ | |
| **Definition of** $(m_1)^{(3)}, (m_2)^{(3)}, (\mu_1)^{(3)}, (\mu_2)^{(3)}$ :-<br>If we define $(m_1)^{(3)}, (m_2)^{(3)}, (\mu_1)^{(3)}, (\mu_2)^{(3)}$ by<br>$(m_2)^{(3)} = (v_0)^{(3)}, (m_1)^{(3)} = (v_1)^{(3)}$, ***if*** $(v_0)^{(3)} < (v_1)^{(3)}$<br>$(m_2)^{(3)} = (v_1)^{(3)}, (m_1)^{(3)} = (\bar{v}_1)^{(3)}$, ***if*** $(v_1)^{(3)} < (v_0)^{(3)} < (\bar{v}_1)^{(3)}$,<br>and $\boxed{(v_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0}}$<br>$(m_2)^{(3)} = (v_1)^{(3)}, (m_1)^{(3)} = (v_0)^{(3)}$, ***if*** $(\bar{v}_1)^{(3)} < (v_0)^{(3)}$ | 321 |
| and analogously<br>$(\mu_2)^{(3)} = (u_0)^{(3)}, (\mu_1)^{(3)} = (u_1)^{(3)}$, ***if*** $(u_0)^{(3)} < (u_1)^{(3)}$<br>$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (\bar{u}_1)^{(3)}$, ***if*** $(u_1)^{(3)} < (u_0)^{(3)} < (\bar{u}_1)^{(3)}$, and $\boxed{(u_0)^{(3)} = \frac{T_{20}^0}{T_{21}^0}}$<br>$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (u_0)^{(3)}$, ***if*** $(\bar{u}_1)^{(3)} < (u_0)^{(3)}$<br>Then the solution of global equations satisfies the inequalities<br>$G_{20}^0 e^{((S_1)^{(3)} - (p_{20})^{(3)})t} \le G_{20}(t) \le G_{20}^0 e^{(S_1)^{(3)}t}$<br>$(p_i)^{(3)}$ is defined by equation | 322 |
| $\frac{1}{(m_1)^{(3)}} G_{20}^0 e^{((S_1)^{(3)} - (p_{20})^{(3)})t} \le G_{21}(t) \le \frac{1}{(m_2)^{(3)}} G_{20}^0 e^{(S_1)^{(3)}t}$ | 323 |
| $\left(\frac{(a_{22})^{(3)} G_{20}^0}{(m_1)^{(3)}((S_1)^{(3)} - (p_{20})^{(3)} - (S_2)^{(3)})}\left[e^{((S_1)^{(3)} - (p_{20})^{(3)})t} - e^{-(S_2)^{(3)}t}\right] + G_{22}^0 e^{-(S_2)^{(3)}t} \le G_{22}(t) \le \right.$<br>$\left.\frac{(a_{22})^{(3)} G_{20}^0}{(m_2)^{(3)}((S_1)^{(3)} - (a'_{22})^{(3)})}\left[e^{(S_1)^{(3)}t} - e^{-(a'_{22})^{(3)}t}\right] + G_{22}^0 e^{-(a'_{22})^{(3)}t}\right)$ | 324 |
| $\boxed{T_{20}^0 e^{(R_1)^{(3)}t} \le T_{20}(t) \le T_{20}^0 e^{((R_1)^{(3)} + (r_{20})^{(3)})t}}$ | 325 |
| $\frac{1}{(\mu_1)^{(3)}} T_{20}^0 e^{(R_1)^{(3)}t} \le T_{20}(t) \le \frac{1}{(\mu_2)^{(3)}} T_{20}^0 e^{((R_1)^{(3)} + (r_{20})^{(3)})t}$ | 326 |
| $\frac{(b_{22})^{(3)} T_{20}^0}{(\mu_1)^{(3)}((R_1)^{(3)} - (b'_{22})^{(3)})}\left[e^{(R_1)^{(3)}t} - e^{-(b'_{22})^{(3)}t}\right] + T_{22}^0 e^{-(b'_{22})^{(3)}t} \le T_{22}(t) \le$<br>$\frac{(a_{22})^{(3)} T_{20}^0}{(\mu_2)^{(3)}((R_1)^{(3)} + (r_{20})^{(3)} + (R_2)^{(3)})}\left[e^{((R_1)^{(3)} + (r_{20})^{(3)})t} - e^{-(R_2)^{(3)}t}\right] + T_{22}^0 e^{-(R_2)^{(3)}t}$ | 327 |

**Definition of** $(S_1)^{(3)}, (S_2)^{(3)}, (R_1)^{(3)}, (R_2)^{(3)}$:- 328

Where $(S_1)^{(3)} = (a_{20})^{(3)}(m_2)^{(3)} - (a'_{20})^{(3)}$

$(S_2)^{(3)} = (a_{22})^{(3)} - (p_{22})^{(3)}$

$(R_1)^{(3)} = (b_{20})^{(3)}(\mu_2)^{(3)} - (b'_{20})^{(3)}$

$(R_2)^{(3)} = (b'_{22})^{(3)} - (r_{22})^{(3)}$

**Behavior of the solutions of equation**

**Theorem:** If we denote and define

**Definition of** $(\sigma_1)^{(4)}, (\sigma_2)^{(4)}, (\tau_1)^{(4)}, (\tau_2)^{(4)}$ :

$(\sigma_1)^{(4)}, (\sigma_2)^{(4)}, (\tau_1)^{(4)}, (\tau_2)^{(4)}$ four constants satisfying

$-(\sigma_2)^{(4)} \le -(a'_{24})^{(4)} + (a'_{25})^{(4)} - (a''_{24})^{(4)}(T_{25}, t) + (a''_{25})^{(4)}(T_{25}, t) \le -(\sigma_1)^{(4)}$

$-(\tau_2)^{(4)} \le -(b'_{24})^{(4)} + (b'_{25})^{(4)} - (b''_{24})^{(4)}((G_{27}), t) - (b''_{25})^{(4)}((G_{27}), t) \le -(\tau_1)^{(4)}$

**Definition of** $(v_1)^{(4)}, (v_2)^{(4)}, (u_1)^{(4)}, (u_2)^{(4)}, v^{(4)}, u^{(4)}$ : 329

By $(v_1)^{(4)} > 0$, $(v_2)^{(4)} < 0$ and respectively $(u_1)^{(4)} > 0$, $(u_2)^{(4)} < 0$ the roots of the equations $(a_{25})^{(4)}\left(v^{(4)}\right)^2 + (\sigma_1)^{(4)}v^{(4)} - (a_{24})^{(4)} = 0$ and $(b_{25})^{(4)}\left(u^{(4)}\right)^2 + (\tau_1)^{(4)}u^{(4)} - (b_{24})^{(4)} = 0$ and

**Definition of** $(\bar{v}_1)^{(4)}, , (\bar{v}_2)^{(4)}, (\bar{u}_1)^{(4)}, (\bar{u}_2)^{(4)}$ : 330

By $(\bar{v}_1)^{(4)} > 0$, $(\bar{v}_2)^{(4)} < 0$ and respectively $(\bar{u}_1)^{(4)} > 0$, $(\bar{u}_2)^{(4)} < 0$ the roots of the equations $(a_{25})^{(4)}\left(v^{(4)}\right)^2 + (\sigma_2)^{(4)}v^{(4)} - (a_{24})^{(4)} = 0$ and $(b_{25})^{(4)}\left(u^{(4)}\right)^2 + (\tau_2)^{(4)}u^{(4)} - (b_{24})^{(4)} = 0$

**Definition of** $(m_1)^{(4)}, (m_2)^{(4)}, (\mu_1)^{(4)}, (\mu_2)^{(4)}, (v_0)^{(4)}$ :-

If we define $(m_1)^{(4)}, (m_2)^{(4)}, (\mu_1)^{(4)}, (\mu_2)^{(4)}$ by

$(m_2)^{(4)} = (v_0)^{(4)}, (m_1)^{(4)} = (v_1)^{(4)}$, ***if*** $(v_0)^{(4)} < (v_1)^{(4)}$

$(m_2)^{(4)} = (v_1)^{(4)}, (m_1)^{(4)} = (\bar{v}_1)^{(4)}$, ***if*** $(v_4)^{(4)} < (v_0)^{(4)} < (\bar{v}_1)^{(4)}$,

and $\boxed{(v_0)^{(4)} = \frac{G^0_{24}}{G^0_{25}}}$

$(m_2)^{(4)} = (v_4)^{(4)}, (m_1)^{(4)} = (v_0)^{(4)}$, ***if*** $(\bar{v}_4)^{(4)} < (v_0)^{(4)}$

and analogously 331

$(\mu_2)^{(4)} = (u_0)^{(4)}, (\mu_1)^{(4)} = (u_1)^{(4)}$, ***if*** $(u_0)^{(4)} < (u_1)^{(4)}$

$(\mu_2)^{(4)} = (u_1)^{(4)}, (\mu_1)^{(4)} = (\bar{u}_1)^{(4)}$, ***if*** $(u_1)^{(4)} < (u_0)^{(4)} < (\bar{u}_1)^{(4)}$,

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| and $\boxed{(u_0)^{(4)} = \frac{T_{24}^0}{T_{25}^0}}$<br>$(\mu_2)^{(4)} = (u_1)^{(4)}, (\mu_1)^{(4)} = (u_0)^{(4)}, \boldsymbol{if}\ (\bar{u}_1)^{(4)} < (u_0)^{(4)}$ where $(u_1)^{(4)}, (\bar{u}_1)^{(4)}$ | |
| Then the solution of global equations satisfies the inequalities<br>$G_{24}^0 e^{((S_1)^{(4)} - (p_{24})^{(4)})t} \le G_{24}(t) \le G_{24}^0 e^{(S_1)^{(4)} t}$<br>where $(p_i)^{(4)}$ is defined by equation | 332 |
| $\frac{1}{(m_1)^{(4)}} G_{24}^0 e^{((S_1)^{(4)} - (p_{24})^{(4)})t} \le G_{25}(t) \le \frac{1}{(m_2)^{(4)}} G_{24}^0 e^{(S_1)^{(4)} t}$ | 333 |
| $\left(\frac{(a_{26})^{(4)} G_{24}^0}{(m_1)^{(4)}((S_1)^{(4)} - (p_{24})^{(4)} - (S_2)^{(4)})}\right)\left[e^{((S_1)^{(4)} - (p_{24})^{(4)})t} - e^{-(S_2)^{(4)} t}\right] + G_{26}^0 e^{-(S_2)^{(4)} t} \le G_{26}(t) \le$<br>*(a26)4G240(m2)4(S1)4−(a26′)4e(S1)4t−e−(a26′)4t+ G260e−(a26′)4t* | 334 |
| $\boxed{T_{24}^0 e^{(R_1)^{(4)} t} \le T_{24}(t) \le T_{24}^0 e^{((R_1)^{(4)} + (r_{24})^{(4)})t}}$ | |
| $\frac{1}{(\mu_1)^{(4)}} T_{24}^0 e^{(R_1)^{(4)} t} \le T_{24}(t) \le \frac{1}{(\mu_2)^{(4)}} T_{24}^0 e^{((R_1)^{(4)} + (r_{24})^{(4)})t}$ | 335 |
| $\frac{(b_{26})^{(4)} T_{24}^0}{(\mu_1)^{(4)}((R_1)^{(4)} - (b_{26}')^{(4)})}\left[e^{(R_1)^{(4)} t} - e^{-(b_{26}')^{(4)} t}\right] + T_{26}^0 e^{-(b_{26}')^{(4)} t} \le T_{26}(t) \le$<br>$\frac{(a_{26})^{(4)} T_{24}^0}{(\mu_2)^{(4)}((R_1)^{(4)} + (r_{24})^{(4)} + (R_2)^{(4)})}\left[e^{((R_1)^{(4)} + (r_{24})^{(4)})t} - e^{-(R_2)^{(4)} t}\right] + T_{26}^0 e^{-(R_2)^{(4)} t}$ | 336 |
| **Definition of** $(S_1)^{(4)}, (S_2)^{(4)}, (R_1)^{(4)}, (R_2)^{(4)}$:-<br>Where $(S_1)^{(4)} = (a_{24})^{(4)}(m_2)^{(4)} - (a_{24}')^{(4)}$<br>$(S_2)^{(4)} = (a_{26})^{(4)} - (p_{26})^{(4)}$<br>$(R_1)^{(4)} = (b_{24})^{(4)}(\mu_2)^{(4)} - (b_{24}')^{(4)}$<br>$(R_2)^{(4)} = (b_{26}')^{(4)} - (r_{26})^{(4)}$ | 337 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ :<br>$(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ four constants satisfying<br>$-(\sigma_2)^{(5)} \le -(a_{28}')^{(5)} + (a_{29}')^{(5)} - (a_{28}'')^{(5)}(T_{29}, t) + (a_{29}'')^{(5)}(T_{29}, t) \le -(\sigma_1)^{(5)}$<br>$-(\tau_2)^{(5)} \le -(b_{28}')^{(5)} + (b_{29}')^{(5)} - (b_{28}'')^{(5)}((G_{31}), t) - (b_{29}'')^{(5)}((G_{31}), t) \le -(\tau_1)^{(5)}$ | 338 |

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| **Definition of** $(v_1)^{(5)}, (v_2)^{(5)}, (u_1)^{(5)}, (u_2)^{(5)}, v^{(5)}, u^{(5)}$ :<br><br>By $(v_1)^{(5)} > 0$ , $(v_2)^{(5)} < 0$ and respectively $(u_1)^{(5)} > 0$ , $(u_2)^{(5)} < 0$ the roots of the equations $(a_{29})^{(5)}\left(v^{(5)}\right)^2 + (\sigma_1)^{(5)}v^{(5)} - (a_{28})^{(5)} = 0$<br>and $(b_{29})^{(5)}\left(u^{(5)}\right)^2 + (\tau_1)^{(5)}u^{(5)} - (b_{28})^{(5)} = 0$ and | 339 |
| **Definition of** $(\bar{v}_1)^{(5)},, (\bar{v}_2)^{(5)}, (\bar{u}_1)^{(5)}, (\bar{u}_2)^{(5)}$ :<br><br>By $(\bar{v}_1)^{(5)} > 0$ , $(\bar{v}_2)^{(5)} < 0$ and respectively $(\bar{u}_1)^{(5)} > 0$ , $(\bar{u}_2)^{(5)} < 0$ the roots of the equations $(a_{29})^{(5)}\left(v^{(5)}\right)^2 + (\sigma_2)^{(5)}v^{(5)} - (a_{28})^{(5)} = 0$<br>and $(b_{29})^{(5)}\left(u^{(5)}\right)^2 + (\tau_2)^{(5)}u^{(5)} - (b_{28})^{(5)} = 0$<br>**Definition of** $(m_1)^{(5)}$ , $(m_2)^{(5)}$ , $(\mu_1)^{(5)}, (\mu_2)^{(5)}, (v_0)^{(5)}$ :-<br><br>If we define $(m_1)^{(5)}$ , $(m_2)^{(5)}$ , $(\mu_1)^{(5)}, (\mu_2)^{(5)}$ by<br><br>$(m_2)^{(5)} = (v_0)^{(5)}, (m_1)^{(5)} = (v_1)^{(5)}$, ***if*** $(v_0)^{(5)} < (v_1)^{(5)}$<br><br>$(m_2)^{(5)} = (v_1)^{(5)}, (m_1)^{(5)} = (\bar{v}_1)^{(5)}$ , ***if*** $(v_1)^{(5)} < (v_0)^{(5)} < (\bar{v}_1)^{(5)}$,<br>and $\boxed{(v_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}}$<br><br>$(m_2)^{(5)} = (v_1)^{(5)}, (m_1)^{(5)} = (v_0)^{(5)}$, ***if*** $(\bar{v}_1)^{(5)} < (v_0)^{(5)}$ | 340 |
| and analogously<br><br>$(\mu_2)^{(5)} = (u_0)^{(5)}, (\mu_1)^{(5)} = (u_1)^{(5)}$, ***if*** $(u_0)^{(5)} < (u_1)^{(5)}$<br><br>$(\mu_2)^{(5)} = (u_1)^{(5)}, (\mu_1)^{(5)} = (\bar{u}_1)^{(5)}$ , ***if*** $(u_1)^{(5)} < (u_0)^{(5)} < (\bar{u}_1)^{(5)}$,<br>and $\boxed{(u_0)^{(5)} = \frac{T_{28}^0}{T_{29}^0}}$<br><br>$(\mu_2)^{(5)} = (u_1)^{(5)}, (\mu_1)^{(5)} = (u_0)^{(5)}$, ***if*** $(\bar{u}_1)^{(5)} < (u_0)^{(5)}$ where $(u_1)^{(5)}, (\bar{u}_1)^{(5)}$ | 341 |
| Then the solution of global equations satisfies the inequalities<br><br>$G_{28}^0 e^{((S_1)^{(5)} - (p_{28})^{(5)})t} \le G_{28}(t) \le G_{28}^0 e^{(S_1)^{(5)}t}$<br><br>where $(p_i)^{(5)}$ is defined by equation | 342 |
| $\frac{1}{(m_5)^{(5)}} G_{28}^0 e^{((S_1)^{(5)} - (p_{28})^{(5)})t} \le G_{29}(t) \le \frac{1}{(m_2)^{(5)}} G_{28}^0 e^{(S_1)^{(5)}t}$ | 343 |
| $\left(\frac{(a_{30})^{(5)} G_{28}^0}{(m_1)^{(5)}((S_1)^{(5)} - (p_{28})^{(5)} - (S_2)^{(5)})}\left[e^{((S_1)^{(5)} - (p_{28})^{(5)})t} - e^{-(S_2)^{(5)}t}\right] + G_{30}^0 e^{-(S_2)^{(5)}t} \le G_{30}(t) \le$<br>*(a30)5G280(m2)5(S1)5−(a30′)5e(S1)5t−e−(a30′)5t+ G300e−(a30′)5t* | 344 |
| $\boxed{T_{28}^0 e^{(R_1)^{(5)}t} \le T_{28}(t) \le T_{28}^0 e^{((R_1)^{(5)} + (r_{28})^{(5)})t}}$ | 345 |

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| $\frac{1}{(\mu_1)^{(5)}}T_{28}^0 e^{(R_1)^{(5)}t} \le T_{28}(t) \le \frac{1}{(\mu_2)^{(5)}}T_{28}^0 e^{((R_1)^{(5)}+(r_{28})^{(5)})t}$ | 346 |
| $\frac{(b_{30})^{(5)}T_{28}^0}{(\mu_1)^{(5)}((R_1)^{(5)}-(b_{30}^{'})^{(5)})}\left[e^{(R_1)^{(5)}t} - e^{-(b_{30}^{'})^{(5)}t}\right] + T_{30}^0 e^{-(b_{30}^{'})^{(5)}t} \le T_{30}(t) \le$ $\frac{(a_{30})^{(5)}T_{28}^0}{(\mu_2)^{(5)}((R_1)^{(5)}+(r_{28})^{(5)}+(R_2)^{(5)})}\left[e^{((R_1)^{(5)}+(r_{28})^{(5)})t} - e^{-(R_2)^{(5)}t}\right] + T_{30}^0 e^{-(R_2)^{(5)}t}$ | 347 |
| **Definition of** $(S_1)^{(5)}, (S_2)^{(5)}, (R_1)^{(5)}, (R_2)^{(5)}$:- Where $(S_1)^{(5)} = (a_{28})^{(5)}(m_2)^{(5)} - (a_{28}^{'})^{(5)}$ $(S_2)^{(5)} = (a_{30})^{(5)} - (p_{30})^{(5)}$ $(R_1)^{(5)} = (b_{28})^{(5)}(\mu_2)^{(5)} - (b_{28}^{'})^{(5)}$ $(R_2)^{(5)} = (b_{30}^{'})^{(5)} - (r_{30})^{(5)}$ | 348 |
| **Behavior of the solutions of equation** **Theorem 2:** If we denote and define **Definition of** $(\sigma_1)^{(6)}, (\sigma_2)^{(6)}, (\tau_1)^{(6)}, (\tau_2)^{(6)}$ : $(\sigma_1)^{(6)}, (\sigma_2)^{(6)}, (\tau_1)^{(6)}, (\tau_2)^{(6)}$ four constants satisfying $-(\sigma_2)^{(6)} \le -(a_{32}^{'})^{(6)} + (a_{33}^{'})^{(6)} - (a_{32}^{''})^{(6)}(T_{33}, t) + (a_{33}^{''})^{(6)}(T_{33}, t) \le -(\sigma_1)^{(6)}$ $-(\tau_2)^{(6)} \le -(b_{32}^{'})^{(6)} + (b_{33}^{'})^{(6)} - (b_{32}^{''})^{(6)}\left((G_{35}), t\right) - (b_{33}^{''})^{(6)}\left((G_{35}), t\right) \le -(\tau_1)^{(6)}$ | 349 |
| **Definition of** $(v_1)^{(6)}, (v_2)^{(6)}, (u_1)^{(6)}, (u_2)^{(6)}, v^{(6)}, u^{(6)}$ : By $(v_1)^{(6)} > 0$, $(v_2)^{(6)} < 0$ and respectively $(u_1)^{(6)} > 0$, $(u_2)^{(6)} < 0$ the roots of the equations $(a_{33})^{(6)}\left(v^{(6)}\right)^2 + (\sigma_1)^{(6)}v^{(6)} - (a_{32})^{(6)} = 0$ and $(b_{33})^{(6)}\left(u^{(6)}\right)^2 + (\tau_1)^{(6)}u^{(6)} - (b_{32})^{(6)} = 0$ and | 350 |
| **Definition of** $(\bar{v}_1)^{(6)},, (\bar{v}_2)^{(6)}, (\bar{u}_1)^{(6)}, (\bar{u}_2)^{(6)}$ : By $(\bar{v}_1)^{(6)} > 0$, $(\bar{v}_2)^{(6)} < 0$ and respectively $(\bar{u}_1)^{(6)} > 0$, $(\bar{u}_2)^{(6)} < 0$ the roots of the equations $(a_{33})^{(6)}\left(v^{(6)}\right)^2 + (\sigma_2)^{(6)}v^{(6)} - (a_{32})^{(6)} = 0$ and $(b_{33})^{(6)}\left(u^{(6)}\right)^2 + (\tau_2)^{(6)}u^{(6)} - (b_{32})^{(6)} = 0$ **Definition of** $(m_1)^{(6)}, (m_2)^{(6)}, (\mu_1)^{(6)}, (\mu_2)^{(6)}, (v_0)^{(6)}$ :- If we define $(m_1)^{(6)}, (m_2)^{(6)}, (\mu_1)^{(6)}, (\mu_2)^{(6)}$ by $(m_2)^{(6)} = (v_0)^{(6)}, (m_1)^{(6)} = (v_1)^{(6)}$, ***if*** $(v_0)^{(6)} < (v_1)^{(6)}$ $(m_2)^{(6)} = (v_1)^{(6)}, (m_1)^{(6)} = (\bar{v}_6)^{(6)}$, ***if*** $(v_1)^{(6)} < (v_0)^{(6)} < (\bar{v}_1)^{(6)}$, | 351 |

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| and $\boxed{(v_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0}}$<br>$(m_2)^{(6)} = (v_1)^{(6)}, (m_1)^{(6)} = (v_0)^{(6)},$ ***if*** $(\bar{v}_1)^{(6)} < (v_0)^{(6)}$ | |
| and analogously<br>$(\mu_2)^{(6)} = (u_0)^{(6)}, (\mu_1)^{(6)} = (u_1)^{(6)},$ ***if*** $(u_0)^{(6)} < (u_1)^{(6)}$<br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (\bar{u}_1)^{(6)}$ , ***if*** $(u_1)^{(6)} < (u_0)^{(6)} < (\bar{u}_1)^{(6)}$,<br>and $\boxed{(u_0)^{(6)} = \frac{T_{32}^0}{T_{33}^0}}$<br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (u_0)^{(6)},$ ***if*** $(\bar{u}_1)^{(6)} < (u_0)^{(6)}$ where $(u_1)^{(6)}, (\bar{u}_1)^{(6)}$ | 352 |
| Then the solution of global equations satisfies the inequalities<br>$G_{32}^0 e^{((S_1)^{(6)} - (p_{32})^{(6)})t} \le G_{32}(t) \le G_{32}^0 e^{(S_1)^{(6)} t}$<br>where $(p_i)^{(6)}$ is defined by equation | 353 |
| $\frac{1}{(m_1)^{(6)}} G_{32}^0 e^{((S_1)^{(6)} - (p_{32})^{(6)})t} \le G_{33}(t) \le \frac{1}{(m_2)^{(6)}} G_{32}^0 e^{(S_1)^{(6)} t}$ | 354 |
| $\left(\frac{(a_{34})^{(6)} G_{32}^0}{(m_1)^{(6)}((S_1)^{(6)} - (p_{32})^{(6)} - (S_2)^{(6)})}\right)\left[e^{((S_1)^{(6)} - (p_{32})^{(6)})t} - e^{-(S_2)^{(6)} t}\right] + G_{34}^0 e^{-(S_2)^{(6)} t} \le G_{34}(t) \le$<br>*(a34)6G320(m2)6(S1)6−(a34')6e(S1)6t−e−(a34')6t+ G340e−(a34')6t* | 355 |
| $\boxed{T_{32}^0 e^{(R_1)^{(6)} t} \le T_{32}(t) \le T_{32}^0 e^{((R_1)^{(6)} + (r_{32})^{(6)})t}}$ | 356 |
| $\frac{1}{(\mu_1)^{(6)}} T_{32}^0 e^{(R_1)^{(6)} t} \le T_{32}(t) \le \frac{1}{(\mu_2)^{(6)}} T_{32}^0 e^{((R_1)^{(6)} + (r_{32})^{(6)})t}$ | 357 |
| $\frac{(b_{34})^{(6)} T_{32}^0}{(\mu_1)^{(6)}((R_1)^{(6)} - (b_{34}')^{(6)})}\left[e^{(R_1)^{(6)} t} - e^{-(b_{34}')^{(6)} t}\right] + T_{34}^0 e^{-(b_{34}')^{(6)} t} \le T_{34}(t) \le$<br>$\frac{(a_{34})^{(6)} T_{32}^0}{(\mu_2)^{(6)}((R_1)^{(6)} + (r_{32})^{(6)} + (R_2)^{(6)})}\left[e^{((R_1)^{(6)} + (r_{32})^{(6)})t} - e^{-(R_2)^{(6)} t}\right] + T_{34}^0 e^{-(R_2)^{(6)} t}$ | 358 |
| **Definition of** $(S_1)^{(6)}, (S_2)^{(6)}, (R_1)^{(6)}, (R_2)^{(6)}$:-<br>Where $(S_1)^{(6)} = (a_{32})^{(6)}(m_2)^{(6)} - (a_{32}')^{(6)}$<br>$(S_2)^{(6)} = (a_{34})^{(6)} - (p_{34})^{(6)}$<br>$(R_1)^{(6)} = (b_{32})^{(6)}(\mu_2)^{(6)} - (b_{32}')^{(6)}$<br>$(R_2)^{(6)} = (b_{34}')^{(6)} - (r_{34})^{(6)}$ | 359 |
| **Behavior of the solutions of equation** | |

**Theorem 2:** If we denote and define

**Definition of** $(\sigma_1)^{(7)}, (\sigma_2)^{(7)}, (\tau_1)^{(7)}, (\tau_2)^{(7)}$ :

$(\sigma_1)^{(7)}, (\sigma_2)^{(7)}, (\tau_1)^{(7)}, (\tau_2)^{(7)}$ four constants satisfying

$$-(\sigma_2)^{(7)} \le -(a'_{36})^{(7)} + (a'_{37})^{(7)} - (a''_{36})^{(7)}(T_{37}, t) + (a''_{37})^{(7)}(T_{37}, t) \le -(\sigma_1)^{(7)}$$

$$-(\tau_2)^{(7)} \le -(b'_{36})^{(7)} + (b'_{37})^{(7)} - (b''_{36})^{(7)}\big((G_{39}), t\big) - (b''_{37})^{(7)}\big((G_{39}), t\big) \le -(\tau_1)^{(7)}$$

**Definition of** $(\nu_1)^{(7)}, (\nu_2)^{(7)}, (u_1)^{(7)}, (u_2)^{(7)}, \nu^{(7)}, u^{(7)}$ : 361

By $(\nu_1)^{(7)} > 0, (\nu_2)^{(7)} < 0$ and respectively $(u_1)^{(7)} > 0, (u_2)^{(7)} < 0$ the roots of the equations $(a_{37})^{(7)}\big(\nu^{(7)}\big)^2 + (\sigma_1)^{(7)}\nu^{(7)} - (a_{36})^{(7)} = 0$
and $(b_{37})^{(7)}\big(u^{(7)}\big)^2 + (\tau_1)^{(7)}u^{(7)} - (b_{36})^{(7)} = 0$ and

**Definition of** $(\bar{\nu}_1)^{(7)}, , (\bar{\nu}_2)^{(7)}, (\bar{u}_1)^{(7)}, (\bar{u}_2)^{(7)}$ : 362

By $(\bar{\nu}_1)^{(7)} > 0, (\bar{\nu}_2)^{(7)} < 0$ and respectively $(\bar{u}_1)^{(7)} > 0, (\bar{u}_2)^{(7)} < 0$ the roots of the equations $(a_{37})^{(7)}\big(\nu^{(7)}\big)^2 + (\sigma_2)^{(7)}\nu^{(7)} - (a_{36})^{(7)} = 0$
and $(b_{37})^{(7)}\big(u^{(7)}\big)^2 + (\tau_2)^{(7)}u^{(7)} - (b_{36})^{(7)} = 0$

**Definition of** $(m_1)^{(7)}, (m_2)^{(7)}, (\mu_1)^{(7)}, (\mu_2)^{(7)}, (\nu_0)^{(7)}$ :-

If we define $(m_1)^{(7)}, (m_2)^{(7)}, (\mu_1)^{(7)}, (\mu_2)^{(7)}$ by

$(m_2)^{(7)} = (\nu_0)^{(7)}, (m_1)^{(7)} = (\nu_1)^{(7)}$, ***if*** $(\nu_0)^{(7)} < (\nu_1)^{(7)}$

$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\bar{\nu}_1)^{(7)}$, ***if*** $(\nu_1)^{(7)} < (\nu_0)^{(7)} < (\bar{\nu}_1)^{(7)}$,
and $\boxed{(\nu_0)^{(7)} = \frac{G^0_{36}}{G^0_{37}}}$

$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\nu_0)^{(7)}$, ***if*** $(\bar{\nu}_1)^{(7)} < (\nu_0)^{(7)}$

and analogously 363

$(\mu_2)^{(7)} = (u_0)^{(7)}, (\mu_1)^{(7)} = (u_1)^{(7)}$, ***if*** $(u_0)^{(7)} < (u_1)^{(7)}$

$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (\bar{u}_1)^{(7)}$, ***if*** $(u_1)^{(7)} < (u_0)^{(7)} < (\bar{u}_1)^{(7)}$,
and $\boxed{(u_0)^{(7)} = \frac{T^0_{36}}{T^0_{37}}}$

$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (u_0)^{(7)}$, ***if*** $(\bar{u}_1)^{(7)} < (u_0)^{(7)}$ where $(u_1)^{(7)}, (\bar{u}_1)^{(7)}$

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| Then the solution of global equations satisfies the inequalities<br>$G_{36}^0 e^{((S_1)^{(7)}-(p_{36})^{(7)})t} \le G_{36}(t) \le G_{36}^0 e^{(S_1)^{(7)}t}$<br>where $(p_i)^{(7)}$ is defined by equation | 364 |
| $\frac{1}{(m_7)^{(7)}} G_{36}^0 e^{((S_1)^{(7)}-(p_{36})^{(7)})t} \le G_{37}(t) \le \frac{1}{(m_2)^{(7)}} G_{36}^0 e^{(S_1)^{(7)}t}$ | 365 |
| $(\frac{(a_{38})^{(7)} G_{36}^0}{(m_1)^{(7)}((S_1)^{(7)}-(p_{36})^{(7)}-(S_2)^{(7)})}\left[e^{((S_1)^{(7)}-(p_{36})^{(7)})t} - e^{-(S_2)^{(7)}t}\right] + G_{38}^0 e^{-(S_2)^{(7)}t} \le G_{38}(t) \le$<br>$\frac{(a_{38})^{(7)} G_{36}^0}{(m_2)^{(7)}((S_1)^{(7)}-(a_{38}')^{(7)})}[e^{(S_1)^{(7)}t} - e^{-(a_{38}')^{(7)}t}] + G_{38}^0 e^{-(a_{38}')^{(7)}t})$ | 366 |
| $\boxed{T_{36}^0 e^{(R_1)^{(7)}t} \le T_{36}(t) \le T_{36}^0 e^{((R_1)^{(7)}+(r_{36})^{(7)})t}}$ | 367 |
| $\frac{1}{(\mu_1)^{(7)}} T_{36}^0 e^{(R_1)^{(7)}t} \le T_{36}(t) \le \frac{1}{(\mu_2)^{(7)}} T_{36}^0 e^{((R_1)^{(7)}+(r_{36})^{(7)})t}$ | 368 |
| $\frac{(b_{38})^{(7)} T_{36}^0}{(\mu_1)^{(7)}((R_1)^{(7)}-(b_{38}')^{(7)})}\left[e^{(R_1)^{(7)}t} - e^{-(b_{38}')^{(7)}t}\right] + T_{38}^0 e^{-(b_{38}')^{(7)}t} \le T_{38}(t) \le$<br>$\frac{(a_{38})^{(7)} T_{36}^0}{(\mu_2)^{(7)}((R_1)^{(7)}+(r_{36})^{(7)}+(R_2)^{(7)})}\left[e^{((R_1)^{(7)}+(r_{36})^{(7)})t} - e^{-(R_2)^{(7)}t}\right] + T_{38}^0 e^{-(R_2)^{(7)}t}$ | 369 |
| **Definition of** $(S_1)^{(7)}, (S_2)^{(7)}, (R_1)^{(7)}, (R_2)^{(7)}$:-<br>Where $(S_1)^{(7)} = (a_{36})^{(7)}(m_2)^{(7)} - (a_{36}')^{(7)}$<br>$(S_2)^{(7)} = (a_{38})^{(7)} - (p_{38})^{(7)}$<br>$(R_1)^{(7)} = (b_{36})^{(7)}(\mu_2)^{(7)} - (b_{36}')^{(7)}$<br>$(R_2)^{(7)} = (b_{38}')^{(7)} - (r_{38})^{(7)}$ | 370 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ :<br>$(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ four constants satisfying<br>$-(\sigma_2)^{(8)} \le -(a_{40}')^{(8)} + (a_{41}')^{(8)} - (a_{40}'')^{(8)}(T_{41}, t) + (a_{41}'')^{(8)}(T_{41}, t) \le -(\sigma_1)^{(8)}$<br>$-(\tau_2)^{(8)} \le -(b_{40}')^{(8)} + (b_{41}')^{(8)} - (b_{40}'')^{(8)}((G_{43}), t) - (b_{41}'')^{(8)}((G_{43}), t) \le -(\tau_1)^{(8)}$ | 371 |

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| **Definition of** $(v_1)^{(8)}, (v_2)^{(8)}, (u_1)^{(8)}, (u_2)^{(8)}, v^{(8)}, u^{(8)}$ : <br><br> By $(v_1)^{(8)} > 0$ , $(v_2)^{(8)} < 0$ and respectively $(u_1)^{(8)} > 0$ , $(u_2)^{(8)} < 0$ the roots of the equations $(a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_1)^{(8)}v^{(8)} - (a_{40})^{(8)} = 0$ <br> and $(b_{41})^{(8)}\left(u^{(8)}\right)^2 + (\tau_1)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$ and | 372 |
| **Definition of** $(\bar{v}_1)^{(8)}$,, $(\bar{v}_2)^{(8)}, (\bar{u}_1)^{(8)}, (\bar{u}_2)^{(8)}$ : <br><br> By $(\bar{v}_1)^{(8)} > 0$ , $(\bar{v}_2)^{(8)} < 0$ and respectively $(\bar{u}_1)^{(8)} > 0$ , $(\bar{u}_2)^{(8)} < 0$ the <br><br> roots of the equations $(a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_2)^{(8)}v^{(8)} - (a_{40})^{(8)} = 0$ <br><br> and $(b_{41})^{(8)}\left(u^{(8)}\right)^2 + (\tau_2)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$ <br><br> **Definition of** $(m_1)^{(8)}$ , $(m_2)^{(8)}$ , $(\mu_1)^{(8)}, (\mu_2)^{(8)}, (v_0)^{(8)}$ :- <br><br> If we define $(m_1)^{(8)}$ , $(m_2)^{(8)}$ , $(\mu_1)^{(8)}, (\mu_2)^{(8)}$ by <br><br> $(m_2)^{(8)} = (v_0)^{(8)}, (m_1)^{(8)} = (v_1)^{(8)}$, ***if*** $(v_0)^{(8)} < (v_1)^{(8)}$ <br><br> $(m_2)^{(8)} = (v_1)^{(8)}, (m_1)^{(8)} = (\bar{v}_1)^{(8)}$ , ***if*** $(v_1)^{(8)} < (v_0)^{(8)} < (\bar{v}_1)^{(8)}$, <br> and $\boxed{(v_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}}$ <br><br> $(m_2)^{(8)} = (v_1)^{(8)}, (m_1)^{(8)} = (v_0)^{(8)}$, ***if*** $(\bar{v}_1)^{(8)} < (v_0)^{(8)}$ | |
| and analogously <br><br> $(\mu_2)^{(8)} = (u_0)^{(8)}, (\mu_1)^{(8)} = (u_1)^{(8)}$, ***if*** $(u_0)^{(8)} < (u_1)^{(8)}$ <br><br> $(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (\bar{u}_1)^{(8)}$ , ***if*** $(u_1)^{(8)} < (u_0)^{(8)} < (\bar{u}_1)^{(8)}$, <br> and $\boxed{(u_0)^{(8)} = \frac{T_{40}^0}{T_{41}^0}}$ <br><br> $(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (u_0)^{(8)}$, ***if*** $(\bar{u}_1)^{(8)} < (u_0)^{(8)}$ where $(u_1)^{(8)}, (\bar{u}_1)^{(8)}$ | 374 |
| Then the solution of global equations satisfies the inequalities <br><br> $G_{40}^0 e^{((S_1)^{(8)} - (p_{40})^{(8)})t} \leq G_{40}(t) \leq G_{40}^0 e^{(S_1)^{(8)}t}$ <br><br> where $(p_i)^{(8)}$ is defined by equation | 375 |

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|---|---|
| $\frac{1}{(m_1)^{(8)}} G_{40}^0 e^{((S_1)^{(8)}-(p_{40})^{(8)})t} \le G_{41}(t) \le \frac{1}{(m_2)^{(8)}} G_{40}^0 e^{(S_1)^{(8)}t}$ | 376 |
| $\left(\frac{(a_{42})^{(8)}G_{40}^0}{(m_1)^{(8)}((S_1)^{(8)}-(p_{40})^{(8)}-(S_2)^{(8)})}\left[e^{((S_1)^{(8)}-(p_{40})^{(8)})t} - e^{-(S_2)^{(8)}t}\right] + G_{42}^0 e^{-(S_2)^{(8)}t} \le G_{42}(t) \le \right.$ $\left.\frac{(a_{42})^{(8)}G_{40}^0}{(m_2)^{(8)}((S_1)^{(8)}-(a'_{42})^{(8)})}[e^{(S_1)^{(8)}t} - e^{-(a'_{42})^{(8)}t}] + G_{42}^0 e^{-(a'_{42})^{(8)}t}\right)$ | 377 |
| $\boxed{T_{40}^0 e^{(R_1)^{(8)}t} \le T_{40}(t) \le T_{40}^0 e^{((R_1)^{(8)}+(r_{40})^{(8)})t}}$ | 378 |
| $\frac{1}{(\mu_1)^{(8)}} T_{40}^0 e^{(R_1)^{(8)}t} \le T_{40}(t) \le \frac{1}{(\mu_2)^{(8)}} T_{40}^0 e^{((R_1)^{(8)}+(r_{40})^{(8)})t}$ | 379 |
| $\frac{(b_{42})^{(8)}T_{40}^0}{(\mu_1)^{(8)}((R_1)^{(8)}-(b'_{42})^{(8)})}\left[e^{(R_1)^{(8)}t} - e^{-(b'_{42})^{(8)}t}\right] + T_{42}^0 e^{-(b'_{42})^{(8)}t} \le T_{42}(t) \le$ $\frac{(a_{42})^{(8)}T_{40}^0}{(\mu_2)^{(8)}((R_1)^{(8)}+(r_{40})^{(8)}+(R_2)^{(8)})}\left[e^{((R_1)^{(8)}+(r_{40})^{(8)})t} - e^{-(R_2)^{(8)}t}\right] + T_{42}^0 e^{-(R_2)^{(8)}t}$ | 380 |
| **Definition of** $(S_1)^{(8)}, (S_2)^{(8)}, (R_1)^{(8)}, (R_2)^{(8)}$:-<br>Where $(S_1)^{(8)} = (a_{40})^{(8)}(m_2)^{(8)} - (a'_{40})^{(8)}$<br>$(S_2)^{(8)} = (a_{42})^{(8)} - (p_{42})^{(8)}$<br>$(R_1)^{(8)} = (b_{40})^{(8)}(\mu_2)^{(8)} - (b'_{40})^{(8)}$<br>$(R_2)^{(8)} = (b'_{42})^{(8)} - (r_{42})^{(8)}$ | 381 |
| **Behavior of the solutions of equation 37 to 92**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(9)}, (\sigma_2)^{(9)}, (\tau_1)^{(9)}, (\tau_2)^{(9)}$ :<br>$(\sigma_1)^{(9)}, (\sigma_2)^{(9)}, (\tau_1)^{(9)}, (\tau_2)^{(9)}$ four constants satisfying<br>$-(\sigma_2)^{(9)} \le -(a'_{44})^{(9)} + (a'_{45})^{(9)} - (a''_{44})^{(9)}(T_{45}, t) + (a''_{45})^{(9)}(T_{45}, t) \le -(\sigma_1)^{(9)}$<br>$-(\tau_2)^{(9)} \le -(b'_{44})^{(9)} + (b'_{45})^{(9)} - (b''_{44})^{(9)}((G_{47}), t) - (b''_{45})^{(9)}((G_{47}), t) \le -(\tau_1)^{(9)}$ | 382 |
| **Definition of** $(\nu_1)^{(9)}, (\nu_2)^{(9)}, (u_1)^{(9)}, (u_2)^{(9)}, \nu^{(9)}, u^{(9)}$ :<br>By $(\nu_1)^{(9)} > 0$ , $(\nu_2)^{(9)} < 0$ and respectively $(u_1)^{(9)} > 0$ , $(u_2)^{(9)} < 0$ the roots of the equations $(a_{45})^{(9)}(\nu^{(9)})^2 + (\sigma_1)^{(9)}\nu^{(9)} - (a_{44})^{(9)} = 0$<br>and $(b_{45})^{(9)}(u^{(9)})^2 + (\tau_1)^{(9)}u^{(9)} - (b_{44})^{(9)} = 0$ and | |
| **Definition of** $(\bar{\nu}_1)^{(9)},, (\bar{\nu}_2)^{(9)}, (\bar{u}_1)^{(9)}, (\bar{u}_2)^{(9)}$ : | |

By $(\bar{v}_1)^{(9)} > 0$ , $(\bar{v}_2)^{(9)} < 0$ and respectively $(\bar{u}_1)^{(9)} > 0$ , $(\bar{u}_2)^{(9)} < 0$ the roots of the equations $(a_{45})^{(9)}\left(v^{(9)}\right)^2 + (\sigma_2)^{(9)}v^{(9)} - (a_{44})^{(9)} = 0$ and $(b_{45})^{(9)}\left(u^{(9)}\right)^2 + (\tau_2)^{(9)}u^{(9)} - (b_{44})^{(9)} = 0$

**Definition of** $(m_1)^{(9)}$ , $(m_2)^{(9)}$ , $(\mu_1)^{(9)}$, $(\mu_2)^{(9)}$, $(v_0)^{(9)}$ :-

If we define $(m_1)^{(9)}$ , $(m_2)^{(9)}$ , $(\mu_1)^{(9)}$, $(\mu_2)^{(9)}$ by

$(m_2)^{(9)} = (v_0)^{(9)}, (m_1)^{(9)} = (v_1)^{(9)}$, ***if*** $(v_0)^{(9)} < (v_1)^{(9)}$

$(m_2)^{(9)} = (v_1)^{(9)}, (m_1)^{(9)} = (\bar{v}_1)^{(9)}$ , ***if*** $(v_1)^{(9)} < (v_0)^{(9)} < (\bar{v}_1)^{(9)}$,

and $\boxed{(v_0)^{(9)} = \frac{G_{44}^0}{G_{45}^0}}$

$(m_2)^{(9)} = (v_1)^{(9)}, (m_1)^{(9)} = (v_0)^{(9)}$, ***if*** $(\bar{v}_1)^{(9)} < (v_0)^{(9)}$

and analogously

$(\mu_2)^{(9)} = (u_0)^{(9)}, (\mu_1)^{(9)} = (u_1)^{(9)}$, ***if*** $(u_0)^{(9)} < (u_1)^{(9)}$

$(\mu_2)^{(9)} = (u_1)^{(9)}, (\mu_1)^{(9)} = (\bar{u}_1)^{(9)}$ , ***if*** $(u_1)^{(9)} < (u_0)^{(9)} < (\bar{u}_1)^{(9)}$,

and $\boxed{(u_0)^{(9)} = \frac{T_{44}^0}{T_{45}^0}}$

$(\mu_2)^{(9)} = (u_1)^{(9)}, (\mu_1)^{(9)} = (u_0)^{(9)}$, ***if*** $(\bar{u}_1)^{(9)} < (u_0)^{(9)}$ where $(u_1)^{(9)}, (\bar{u}_1)^{(9)}$ are defined by 59 and 69 respectively

Then the solution of 19,20,21,22,23 and 24 satisfies the inequalities

$$G_{44}^0 e^{((S_1)^{(9)} - (p_{44})^{(9)})t} \le G_{44}(t) \le G_{44}^0 e^{(S_1)^{(9)}t}$$

where $(p_i)^{(9)}$ is defined by equation 45

$$\frac{1}{(m_9)^{(9)}} G_{44}^0 e^{((S_1)^{(9)} - (p_{44})^{(9)})t} \le G_{45}(t) \le \frac{1}{(m_2)^{(9)}} G_{44}^0 e^{(S_1)^{(9)}t}$$

$$\left(\frac{(a_{46})^{(9)} G_{44}^0}{(m_1)^{(9)}((S_1)^{(9)} - (p_{44})^{(9)} - (S_2)^{(9)})}\left[e^{((S_1)^{(9)} - (p_{44})^{(9)})t} - e^{-(S_2)^{(9)}t}\right] + G_{46}^0 e^{-(S_2)^{(9)}t} \le G_{46}(t) \le \right.$$

$$\left. \frac{(a_{46})^{(9)} G_{44}^0}{(m_2)^{(9)}((S_1)^{(9)} - (a'_{46})^{(9)})}\left[e^{(S_1)^{(9)}t} - e^{-(a'_{46})^{(9)}t}\right] + G_{46}^0 e^{-(a'_{46})^{(9)}t}\right)$$

$$\boxed{T_{44}^0 e^{(R_1)^{(9)}t} \le T_{44}(t) \le T_{44}^0 e^{((R_1)^{(9)} + (r_{44})^{(9)})t}}$$

$$\frac{1}{(\mu_1)^{(9)}} T_{44}^0 e^{(R_1)^{(9)}t} \le T_{44}(t) \le \frac{1}{(\mu_2)^{(9)}} T_{44}^0 e^{((R_1)^{(9)} + (r_{44})^{(9)})t}$$

$$\frac{(b_{46})^{(9)} T_{44}^0}{(\mu_1)^{(9)}((R_1)^{(9)} - (b'_{46})^{(9)})}\left[e^{(R_1)^{(9)}t} - e^{-(b'_{46})^{(9)}t}\right] + T_{46}^0 e^{-(b'_{46})^{(9)}t} \le T_{46}(t) \le$$

$$\frac{(a_{46})^{(9)}T_{44}^{0}}{(\mu_2)^{(9)}((R_1)^{(9)}+(r_{44})^{(9)}+(R_2)^{(9)})}\left[e^{((R_1)^{(9)}+(r_{44})^{(9)})t}-e^{-(R_2)^{(9)}t}\right]+T_{46}^{0}e^{-(R_2)^{(9)}t}$$

**Definition of** $(S_1)^{(9)}, (S_2)^{(9)}, (R_1)^{(9)}, (R_2)^{(9)}$:-

Where $(S_1)^{(9)} = (a_{44})^{(9)}(m_2)^{(9)} - (a'_{44})^{(9)}$

$(S_2)^{(9)} = (a_{46})^{(9)} - (p_{46})^{(9)}$

$(R_1)^{(9)} = (b_{44})^{(9)}(\mu_2)^{(9)} - (b'_{44})^{(9)}$

$(R_2)^{(9)} = (b'_{46})^{(9)} - (r_{46})^{(9)}$

**Proof :** From global equations we obtain (383)

$$\frac{dv^{(1)}}{dt} = (a_{13})^{(1)} - \left((a'_{13})^{(1)} - (a'_{14})^{(1)} + (a''_{13})^{(1)}(T_{14},t)\right) - (a''_{14})^{(1)}(T_{14},t)v^{(1)} - (a_{14})^{(1)}v^{(1)}$$

**Definition of** $v^{(1)}$ :- $\boxed{v^{(1)} = \frac{G_{13}}{G_{14}}}$

It follows

$$-\left((a_{14})^{(1)}\left(v^{(1)}\right)^2 + (\sigma_2)^{(1)}v^{(1)} - (a_{13})^{(1)}\right) \le \frac{dv^{(1)}}{dt} \le -\left((a_{14})^{(1)}\left(v^{(1)}\right)^2 + (\sigma_1)^{(1)}v^{(1)} - (a_{13})^{(1)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(1)}, (\nu_0)^{(1)}$ :-

For $0 < \boxed{(\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}} < (\nu_1)^{(1)} < (\bar{\nu}_1)^{(1)}$

$$\nu^{(1)}(t) \ge \frac{(\nu_1)^{(1)}+(C)^{(1)}(\nu_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)}-(\nu_0)^{(1)}\right)t\right]}}{1+(C)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)}-(\nu_0)^{(1)}\right)t\right]}} \quad , \quad \boxed{(C)^{(1)} = \frac{(\nu_1)^{(1)}-(\nu_0)^{(1)}}{(\nu_0)^{(1)}-(\nu_2)^{(1)}}}$$

it follows $(\nu_0)^{(1)} \le \nu^{(1)}(t) \le (\nu_1)^{(1)}$

In the same manner , we get (384)

$$\nu^{(1)}(t) \le \frac{(\bar{\nu}_1)^{(1)}+(\bar{C})^{(1)}(\bar{\nu}_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}}{1+(\bar{C})^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(1)} = \frac{(\bar{\nu}_1)^{(1)}-(\nu_0)^{(1)}}{(\nu_0)^{(1)}-(\bar{\nu}_2)^{(1)}}}$$

From which we deduce $(\nu_0)^{(1)} \le \nu^{(1)}(t) \le (\bar{\nu}_1)^{(1)}$

If $0 < (\nu_1)^{(1)} < (\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0} < (\bar{\nu}_1)^{(1)}$ we find like in the previous case, (385)

$$(\nu_1)^{(1)} \le \frac{(\nu_1)^{(1)}+(C)^{(1)}(\nu_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)}-(\nu_2)^{(1)}\right)t\right]}}{1+(C)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)}-(\nu_2)^{(1)}\right)t\right]}} \le \nu^{(1)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(1)}+(\bar{C})^{(1)}(\bar{\nu}_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}}{1+(\bar{C})^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}} \le (\bar{\nu}_1)^{(1)}$$

If $0 < (\nu_1)^{(1)} \le (\bar{\nu}_1)^{(1)} \le \boxed{(\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}}$ , we obtain 386

$$(\nu_1)^{(1)} \le \nu^{(1)}(t) \le \frac{(\bar{\nu}_1)^{(1)}+(\bar{C})^{(1)}(\bar{\nu}_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}}{1+(\bar{C})^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}} \le (\nu_0)^{(1)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(1)}(t)$ :-

$$(m_2)^{(1)} \le \nu^{(1)}(t) \le (m_1)^{(1)}, \quad \boxed{\nu^{(1)}(t) = \frac{G_{13}(t)}{G_{14}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(1)}(t)$ :-

$$(\mu_2)^{(1)} \le u^{(1)}(t) \le (\mu_1)^{(1)}, \quad \boxed{u^{(1)}(t) = \frac{T_{13}(t)}{T_{14}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{13}^{''})^{(1)} = (a_{14}^{''})^{(1)}, then\ (\sigma_1)^{(1)} = (\sigma_2)^{(1)}$ and in this case $(\nu_1)^{(1)} = (\bar{\nu}_1)^{(1)}$ if in addition $(\nu_0)^{(1)} = (\nu_1)^{(1)}$ then $\nu^{(1)}(t) = (\nu_0)^{(1)}$ and as a consequence $G_{13}(t) = (\nu_0)^{(1)}G_{14}(t)$ this also defines $(\nu_0)^{(1)}$ for the special case

Analogously if $(b_{13}^{''})^{(1)} = (b_{14}^{''})^{(1)}, then\ (\tau_1)^{(1)} = (\tau_2)^{(1)}$ and then

$(u_1)^{(1)} = (\bar{u}_1)^{(1)}$if in addition $(u_0)^{(1)} = (u_1)^{(1)}$ then $T_{13}(t) = (u_0)^{(1)}T_{14}(t)$ This is an important consequence of the relation between $(\nu_1)^{(1)}$ and $(\bar{\nu}_1)^{(1)}$, and definition of $(u_0)^{(1)}$.

**Proof** : From global equations we obtain 387

$$\frac{d\nu^{(2)}}{dt} = (a_{16})^{(2)} - \left((a_{16}')^{(2)} - (a_{17}')^{(2)} + (a_{16}'')^{(2)}(T_{17}, t)\right) - (a_{17}'')^{(2)}(T_{17}, t)\nu^{(2)} - (a_{17})^{(2)}\nu^{(2)}$$

**Definition of** $\nu^{(2)}$ :- $\boxed{\nu^{(2)} = \frac{G_{16}}{G_{17}}}$ 388

It follows 389

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| $-\left((a_{17})^{(2)}\left(v^{(2)}\right)^2+(\sigma_2)^{(2)}v^{(2)}-(a_{16})^{(2)}\right)\le\frac{\mathrm{d}v^{(2)}}{\mathrm{dt}}\le-\left((a_{17})^{(2)}\left(v^{(2)}\right)^2+(\sigma_1)^{(2)}v^{(2)}-(a_{16})^{(2)}\right)$ | |
| From which one obtains<br><br>**Definition of** $(\bar{v}_1)^{(2)},(v_0)^{(2)}$ :-<br><br>For $0<(v_0)^{(2)}=\frac{G_{16}^0}{G_{17}^0}<(v_1)^{(2)}<(\bar{v}_1)^{(2)}$<br><br>$v^{(2)}(t)\ge\frac{(v_1)^{(2)}+(C)^{(2)}(v_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((v_1)^{(2)}-(v_0)^{(2)}\right)t\right]}}{1+(C)^{(2)}e^{\left[-(a_{17})^{(2)}\left((v_1)^{(2)}-(v_0)^{(2)}\right)t\right]}}$ , $\boxed{(C)^{(2)}=\frac{(v_1)^{(2)}-(v_0)^{(2)}}{(v_0)^{(2)}-(v_2)^{(2)}}}$<br><br>it follows $(v_0)^{(2)}\le v^{(2)}(t)\le(v_1)^{(2)}$ | 390 |
| In the same manner , we get<br><br>$v^{(2)}(t)\le\frac{(\bar{v}_1)^{(2)}+(\bar{C})^{(2)}(\bar{v}_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}{1+(\bar{C})^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}$ , $\boxed{(\bar{C})^{(2)}=\frac{(\bar{v}_1)^{(2)}-(v_0)^{(2)}}{(v_0)^{(2)}-(\bar{v}_2)^{(2)}}}$ | 391 |
| From which we deduce $(v_0)^{(2)}\le v^{(2)}(t)\le(\bar{v}_1)^{(2)}$ | 392 |
| If $0<(v_1)^{(2)}<(v_0)^{(2)}=\frac{G_{16}^0}{G_{17}^0}<(\bar{v}_1)^{(2)}$ we find like in the previous case,<br><br>$(v_1)^{(2)}\le\frac{(v_1)^{(2)}+(C)^{(2)}(v_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((v_1)^{(2)}-(v_2)^{(2)}\right)t\right]}}{1+(C)^{(2)}e^{\left[-(a_{17})^{(2)}\left((v_1)^{(2)}-(v_2)^{(2)}\right)t\right]}}\le v^{(2)}(t)\le$<br><br>$\frac{(\bar{v}_1)^{(2)}+(\bar{C})^{(2)}(\bar{v}_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}{1+(\bar{C})^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}\le(\bar{v}_1)^{(2)}$ | 393 |
| If $0<(v_1)^{(2)}\le(\bar{v}_1)^{(2)}\le(v_0)^{(2)}=\frac{G_{16}^0}{G_{17}^0}$ , we obtain<br><br>$(v_1)^{(2)}\le v^{(2)}(t)\le\frac{(\bar{v}_1)^{(2)}+(\bar{C})^{(2)}(\bar{v}_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}{1+(\bar{C})^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}\le(v_0)^{(2)}$<br><br>And so with the notation of the first part of condition (c) , we have | 394 |
| **Definition of** $v^{(2)}(t)$ :-<br><br>$(m_2)^{(2)}\le v^{(2)}(t)\le(m_1)^{(2)}$, $\boxed{v^{(2)}(t)=\frac{G_{16}(t)}{G_{17}(t)}}$ | 395 |
| In a completely analogous way, we obtain<br><br>**Definition of** $u^{(2)}(t)$ :- | 396 |

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| $(\mu_2)^{(2)} \leq u^{(2)}(t) \leq (\mu_1)^{(2)}$, $\boxed{u^{(2)}(t) = \frac{T_{16}(t)}{T_{17}(t)}}$ | |
| Now, using this result and replacing it in global equations we get easily the result stated in the theorem. | |
| **Particular case :**<br><br>If $(a_{16}^{''})^{(2)} = (a_{17}^{''})^{(2)}, then\ (\sigma_1)^{(2)} = (\sigma_2)^{(2)}$ and in this case $(\nu_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ if in addition $(\nu_0)^{(2)} = (\nu_1)^{(2)}$ then $\nu^{(2)}(t) = (\nu_0)^{(2)}$ and as a consequence $G_{16}(t) = (\nu_0)^{(2)} G_{17}(t)$<br><br>Analogously if $(b_{16}^{''})^{(2)} = (b_{17}^{''})^{(2)}, then\ (\tau_1)^{(2)} = (\tau_2)^{(2)}$ and then<br><br>$(u_1)^{(2)} = (\bar{u}_1)^{(2)}$if in addition $(u_0)^{(2)} = (u_1)^{(2)}$ then $T_{16}(t) = (u_0)^{(2)} T_{17}(t)$ This is an important consequence of the relation between $(\nu_1)^{(2)}$ and $(\bar{\nu}_1)^{(2)}$ | 397 |
| **Proof** : From global equations we obtain<br><br>$\frac{d\nu^{(3)}}{dt} = (a_{20})^{(3)} - \left((a_{20}^{'})^{(3)} - (a_{21}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21}, t)\right) - (a_{21}^{''})^{(3)}(T_{21}, t)\nu^{(3)} - (a_{21})^{(3)}\nu^{(3)}$ | 398 |
| **Definition of** $\nu^{(3)}$ :- $\boxed{\nu^{(3)} = \frac{G_{20}}{G_{21}}}$<br><br>It follows<br><br>$-\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_2)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right) \leq \frac{d\nu^{(3)}}{dt} \leq -\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_1)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right)$ | 399 |
| From which one obtains<br><br>For $0 < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\nu_1)^{(3)} < (\bar{\nu}_1)^{(3)}$<br><br>$\nu^{(3)}(t) \geq \frac{(\nu_1)^{(3)} + (C)^{(3)}(\nu_2)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)} - (\nu_0)^{(3)}\right)t\right]}}{1 + (C)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)} - (\nu_0)^{(3)}\right)t\right]}}$ , $\boxed{(C)^{(3)} = \frac{(\nu_1)^{(3)} - (\nu_0)^{(3)}}{(\nu_0)^{(3)} - (\nu_2)^{(3)}}}$<br><br>it follows $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\nu_1)^{(3)}$ | 400 |
| In the same manner , we get<br><br>$\nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)} + (C)^{(3)}(\bar{\nu}_2)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)} - (\bar{\nu}_2)^{(3)}\right)t\right]}}{1 + (C)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)} - (\bar{\nu}_2)^{(3)}\right)t\right]}}$ , $\boxed{(\bar{C})^{(3)} = \frac{(\bar{\nu}_1)^{(3)} - (\nu_0)^{(3)}}{(\nu_0)^{(3)} - (\bar{\nu}_2)^{(3)}}}$<br><br>**Definition of** $(\bar{\nu}_1)^{(3)}$ :-<br><br>From which we deduce $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\bar{\nu}_1)^{(3)}$ | 401 |

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If $0 < (\nu_1)^{(3)} < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\bar{\nu}_1)^{(3)}$ we find like in the previous case,

$$(\nu_1)^{(3)} \leq \frac{(\nu_1)^{(3)}+(C)^{(3)}(\nu_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}}{1+(C)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}} \leq \nu^{(3)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(C)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\bar{\nu}_1)^{(3)}$$

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If $0 < (\nu_1)^{(3)} \leq (\bar{\nu}_1)^{(3)} \leq (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0}$ , we obtain

$$(\nu_1)^{(3)} \leq \nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\nu_0)^{(3)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(3)}(t)$ :-

$$(m_2)^{(3)} \leq \nu^{(3)}(t) \leq (m_1)^{(3)}, \quad \boxed{\nu^{(3)}(t) = \frac{G_{20}(t)}{G_{21}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(3)}(t)$ :-

$$(\mu_2)^{(3)} \leq u^{(3)}(t) \leq (\mu_1)^{(3)}, \quad \boxed{u^{(3)}(t) = \frac{T_{20}(t)}{T_{21}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{20}'')^{(3)} = (a_{21}'')^{(3)}, then\ (\sigma_1)^{(3)} = (\sigma_2)^{(3)}$ and in this case $(\nu_1)^{(3)} = (\bar{\nu}_1)^{(3)}$ if in addition $(\nu_0)^{(3)} = (\nu_1)^{(3)}$ then $\nu^{(3)}(t) = (\nu_0)^{(3)}$ and as a consequence $G_{20}(t) = (\nu_0)^{(3)}G_{21}(t)$

Analogously if $(b_{20}'')^{(3)} = (b_{21}'')^{(3)}, then\ (\tau_1)^{(3)} = (\tau_2)^{(3)}$ and then

$(u_1)^{(3)} = (\bar{u}_1)^{(3)}$if in addition $(u_0)^{(3)} = (u_1)^{(3)}$ then $T_{20}(t) = (u_0)^{(3)}T_{21}(t)$ This is an important consequence of the relation between $(\nu_1)^{(3)}$ and $(\bar{\nu}_1)^{(3)}$

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(4)}}{dt} = (a_{24})^{(4)} - \left((a_{24}')^{(4)} - (a_{25}')^{(4)} + (a_{24}'')^{(4)}(T_{25},t)\right) - (a_{25}'')^{(4)}(T_{25},t)\nu^{(4)} - (a_{25})^{(4)}\nu^{(4)}$$

**Definition of** $\nu^{(4)}$ :- $\boxed{\nu^{(4)} = \frac{G_{24}}{G_{25}}}$

It follows

$$-\left((a_{25})^{(4)}\left(v^{(4)}\right)^2+(\sigma_2)^{(4)}v^{(4)}-(a_{24})^{(4)}\right)\le\frac{dv^{(4)}}{dt}\le-\left((a_{25})^{(4)}\left(v^{(4)}\right)^2+(\sigma_4)^{(4)}v^{(4)}-(a_{24})^{(4)}\right)$$

From which one obtains

**Definition of** $(\bar{v}_1)^{(4)},(v_0)^{(4)}$ :-

For $0<\boxed{(v_0)^{(4)}=\frac{G_{24}^0}{G_{25}^0}}<(v_1)^{(4)}<(\bar{v}_1)^{(4)}$

$$v^{(4)}(t)\ge\frac{(v_1)^{(4)}+(C)^{(4)}(v_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((v_1)^{(4)}-(v_0)^{(4)}\right)t\right]}}{4+(C)^{(4)}e^{\left[-(a_{25})^{(4)}\left((v_1)^{(4)}-(v_0)^{(4)}\right)t\right]}}\quad,\quad\boxed{(C)^{(4)}=\frac{(v_1)^{(4)}-(v_0)^{(4)}}{(v_0)^{(4)}-(v_2)^{(4)}}}$$

it follows $(v_0)^{(4)}\le v^{(4)}(t)\le(v_1)^{(4)}$

In the same manner , we get

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$$v^{(4)}(t)\le\frac{(\bar{v}_1)^{(4)}+(\bar{C})^{(4)}(\bar{v}_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}{4+(\bar{C})^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}\quad,\quad\boxed{(\bar{C})^{(4)}=\frac{(\bar{v}_1)^{(4)}-(v_0)^{(4)}}{(v_0)^{(4)}-(\bar{v}_2)^{(4)}}}$$

From which we deduce $(v_0)^{(4)}\le v^{(4)}(t)\le(\bar{v}_1)^{(4)}$

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If $0<(v_1)^{(4)}<(v_0)^{(4)}=\frac{G_{24}^0}{G_{25}^0}<(\bar{v}_1)^{(4)}$ we find like in the previous case,

$$(v_1)^{(4)}\le\frac{(v_1)^{(4)}+(C)^{(4)}(v_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((v_1)^{(4)}-(v_2)^{(4)}\right)t\right]}}{1+(C)^{(4)}e^{\left[-(a_{25})^{(4)}\left((v_1)^{(4)}-(v_2)^{(4)}\right)t\right]}}\le v^{(4)}(t)\le$$

$$\frac{(\bar{v}_1)^{(4)}+(\bar{C})^{(4)}(\bar{v}_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}{1+(\bar{C})^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}\le(\bar{v}_1)^{(4)}$$

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If $0<(v_1)^{(4)}\le(\bar{v}_1)^{(4)}\le\boxed{(v_0)^{(4)}=\frac{G_{24}^0}{G_{25}^0}}$ , we obtain

$$(v_1)^{(4)}\le v^{(4)}(t)\le\frac{(\bar{v}_1)^{(4)}+(\bar{C})^{(4)}(\bar{v}_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}{1+(\bar{C})^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}\le(v_0)^{(4)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $v^{(4)}(t)$ :-

$$(m_2)^{(4)}\le v^{(4)}(t)\le(m_1)^{(4)},\quad\boxed{v^{(4)}(t)=\frac{G_{24}(t)}{G_{25}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(4)}(t)$ :-

$$(\mu_2)^{(4)}\le u^{(4)}(t)\le(\mu_1)^{(4)},\quad\boxed{u^{(4)}(t)=\frac{T_{24}(t)}{T_{25}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{24}^{''})^{(4)} = (a_{25}^{''})^{(4)}, then\ (\sigma_1)^{(4)} = (\sigma_2)^{(4)}$ and in this case $(\nu_1)^{(4)} = (\bar{\nu}_1)^{(4)}$ if in addition $(\nu_0)^{(4)} = (\nu_1)^{(4)}$ then $\nu^{(4)}(t) = (\nu_0)^{(4)}$ and as a consequence $G_{24}(t) = (\nu_0)^{(4)} G_{25}(t)$ **this also defines** $(\nu_0)^{(4)}$ **for the special case** .

Analogously if $(b_{24}^{''})^{(4)} = (b_{25}^{''})^{(4)}, then\ (\tau_1)^{(4)} = (\tau_2)^{(4)}$ and then $(u_1)^{(4)} = (\bar{u}_4)^{(4)}$ if in addition $(u_0)^{(4)} = (u_1)^{(4)}$ then $T_{24}(t) = (u_0)^{(4)} T_{25}(t)$ This is an important consequence of the relation between $(\nu_1)^{(4)}$ and $(\bar{\nu}_1)^{(4)}$, **and definition of** $(u_0)^{(4)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(5)}}{dt} = (a_{28})^{(5)} - \left((a_{28}^{'})^{(5)} - (a_{29}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}, t)\right) - (a_{29}^{''})^{(5)}(T_{29}, t)\nu^{(5)} - (a_{29})^{(5)}\nu^{(5)}$$

**Definition of** $\nu^{(5)}$ :- $\boxed{\nu^{(5)} = \frac{G_{28}}{G_{29}}}$

It follows

$$-\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_2)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right) \leq \frac{d\nu^{(5)}}{dt} \leq -\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_1)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(5)}, (\nu_0)^{(5)}$ :-

For $0 < \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}} < (\nu_1)^{(5)} < (\bar{\nu}_1)^{(5)}$

$$\nu^{(5)}(t) \geq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}}{5 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}} \quad , \quad \boxed{(C)^{(5)} = \frac{(\nu_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\nu_2)^{(5)}}}$$

it follows $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\nu_1)^{(5)}$

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In the same manner , we get

$$\nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)} + (\bar{C})^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{5 + (\bar{C})^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(5)} = \frac{(\bar{\nu}_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\bar{\nu}_2)^{(5)}}}$$

From which we deduce $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\bar{\nu}_5)^{(5)}$

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If $0 < (\nu_1)^{(5)} < (\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0} < (\bar{\nu}_1)^{(5)}$ we find like in the previous case,

$$(\nu_1)^{(5)} \leq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}}{1 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}} \leq \nu^{(5)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(5)}+(\bar{C})^{(5)}(\bar{\nu}_2)^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}}{1+(\bar{C})^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\bar{\nu}_1)^{(5)}$$

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If $0 < (\nu_1)^{(5)} \leq (\bar{\nu}_1)^{(5)} \leq \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}}$ , we obtain

$$(\nu_1)^{(5)} \leq \nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)}+(\bar{C})^{(5)}(\bar{\nu}_2)^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}}{1+(\bar{C})^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\nu_0)^{(5)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(5)}(t)$ :-

$$(m_2)^{(5)} \leq \nu^{(5)}(t) \leq (m_1)^{(5)}, \quad \boxed{\nu^{(5)}(t) = \frac{G_{28}(t)}{G_{29}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(5)}(t)$ :-

$$(\mu_2)^{(5)} \leq u^{(5)}(t) \leq (\mu_1)^{(5)}, \quad \boxed{u^{(5)}(t) = \frac{T_{28}(t)}{T_{29}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{28}^{''})^{(5)} = (a_{29}^{''})^{(5)}, then\ (\sigma_1)^{(5)} = (\sigma_2)^{(5)}$ and in this case $(\nu_1)^{(5)} = (\bar{\nu}_1)^{(5)}$ if in addition $(\nu_0)^{(5)} = (\nu_5)^{(5)}$ then $\nu^{(5)}(t) = (\nu_0)^{(5)}$ and as a consequence $G_{28}(t) = (\nu_0)^{(5)}G_{29}(t)$ **this also defines** $(\nu_0)^{(5)}$ **for the special case** .

Analogously if $(b_{28}^{''})^{(5)} = (b_{29}^{''})^{(5)}, then\ (\tau_1)^{(5)} = (\tau_2)^{(5)}$ and then $(u_1)^{(5)} = (\bar{u}_1)^{(5)}$ if in addition $(u_0)^{(5)} = (u_1)^{(5)}$ then $T_{28}(t) = (u_0)^{(5)}T_{29}(t)$ This is an important consequence of the relation between $(\nu_1)^{(5)}$ and $(\bar{\nu}_1)^{(5)}$, **and definition of** $(u_0)^{(5)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(6)}}{dt} = (a_{32})^{(6)} - \left((a_{32}^{'})^{(6)} - (a_{33}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33}, t)\right) - (a_{33}^{''})^{(6)}(T_{33}, t)\nu^{(6)} - (a_{33})^{(6)}\nu^{(6)}$$

**Definition of** $\nu^{(6)}$ :- $\boxed{\nu^{(6)} = \frac{G_{32}}{G_{33}}}$

It follows

$$-\left((a_{33})^{(6)}\left(\nu^{(6)}\right)^2 + (\sigma_2)^{(6)}\nu^{(6)} - (a_{32})^{(6)}\right) \leq \frac{d\nu^{(6)}}{dt} \leq -\left((a_{33})^{(6)}\left(\nu^{(6)}\right)^2 + (\sigma_1)^{(6)}\nu^{(6)} - (a_{32})^{(6)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(6)}, (\nu_0)^{(6)}$ :-

For $0 < \boxed{(\nu_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0}} < (\nu_1)^{(6)} < (\bar{\nu}_1)^{(6)}$

$$\nu^{(6)}(t) \geq \frac{(\nu_1)^{(6)}+(C)^{(6)}(\nu_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_0)^{(6)}\right)t\right]}}{1+(C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_0)^{(6)}\right)t\right]}} \quad , \quad \boxed{(C)^{(6)}=\frac{(\nu_1)^{(6)}-(\nu_0)^{(6)}}{(\nu_0)^{(6)}-(\nu_2)^{(6)}}}$$

it follows $(\nu_0)^{(6)} \leq \nu^{(6)}(t) \leq (\nu_1)^{(6)}$

In the same manner , we get

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$$\nu^{(6)}(t) \leq \frac{(\bar{\nu}_1)^{(6)}+(\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}}{1+(\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(6)}=\frac{(\bar{\nu}_1)^{(6)}-(\nu_0)^{(6)}}{(\nu_0)^{(6)}-(\bar{\nu}_2)^{(6)}}}$$

From which we deduce $(\nu_0)^{(6)} \leq \nu^{(6)}(t) \leq (\bar{\nu}_1)^{(6)}$

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If $0 < (\nu_1)^{(6)} < (\nu_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0} < (\bar{\nu}_1)^{(6)}$ we find like in the previous case,

$$(\nu_1)^{(6)} \leq \frac{(\nu_1)^{(6)}+(C)^{(6)}(\nu_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_2)^{(6)}\right)t\right]}}{1+(C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_2)^{(6)}\right)t\right]}} \leq \nu^{(6)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(6)}+(\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}}{1+(\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}} \leq (\bar{\nu}_1)^{(6)}$$

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If $0 < (\nu_1)^{(6)} \leq (\bar{\nu}_1)^{(6)} \leq \boxed{(\nu_0)^{(6)}=\frac{G_{32}^0}{G_{33}^0}}$ , we obtain

$$(\nu_1)^{(6)} \leq \nu^{(6)}(t) \leq \frac{(\bar{\nu}_1)^{(6)}+(\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}}{1+(\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}} \leq (\nu_0)^{(6)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(6)}(t)$ :-

$$(m_2)^{(6)} \leq \nu^{(6)}(t) \leq (m_1)^{(6)}, \quad \boxed{\nu^{(6)}(t)=\frac{G_{32}(t)}{G_{33}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(6)}(t)$ :-

$$(\mu_2)^{(6)} \leq u^{(6)}(t) \leq (\mu_1)^{(6)}, \quad \boxed{u^{(6)}(t)=\frac{T_{32}(t)}{T_{33}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{32}^{''})^{(6)} = (a_{33}^{''})^{(6)}, then\ (\sigma_1)^{(6)} = (\sigma_2)^{(6)}$ and in this case $(\nu_1)^{(6)} = (\bar{\nu}_1)^{(6)}$ if in addition $(\nu_0)^{(6)} = (\nu_1)^{(6)}$ then $\nu^{(6)}(t) = (\nu_0)^{(6)}$ and as a consequence $G_{32}(t) = (\nu_0)^{(6)}G_{33}(t)$ **this also defines** $(\nu_0)^{(6)}$ **for the special case** .

Analogously if $(b_{32}^{''})^{(6)} = (b_{33}^{''})^{(6)}, then\ (\tau_1)^{(6)} = (\tau_2)^{(6)}$ and then $(u_1)^{(6)} = (\bar{u}_1)^{(6)}$ if in addition $(u_0)^{(6)} = (u_1)^{(6)}$ then $T_{32}(t) = (u_0)^{(6)}T_{33}(t)$ This is an important consequence of the relation between $(\nu_1)^{(6)}$ and $(\bar{\nu}_1)^{(6)}$, **and definition of** $(u_0)^{(6)}$.

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**Proof** : From global equations we obtain

$$\frac{dv^{(7)}}{dt} = (a_{36})^{(7)} - \left((a'_{36})^{(7)} - (a'_{37})^{(7)} + (a''_{36})^{(7)}(T_{37},t)\right) - (a''_{37})^{(7)}(T_{37},t)v^{(7)} - (a_{37})^{(7)}v^{(7)}$$

**Definition of** $v^{(7)}$ :- $\boxed{v^{(7)} = \frac{G_{36}}{G_{37}}}$

It follows

$$-\left((a_{37})^{(7)}\left(v^{(7)}\right)^2 + (\sigma_2)^{(7)}v^{(7)} - (a_{36})^{(7)}\right) \le \frac{dv^{(7)}}{dt} \le -\left((a_{37})^{(7)}\left(v^{(7)}\right)^2 + (\sigma_1)^{(7)}v^{(7)} - (a_{36})^{(7)}\right)$$

From which one obtains

**Definition of** $(\bar{v}_1)^{(7)}, (v_0)^{(7)}$ :-

For $0 < \boxed{(v_0)^{(7)} = \frac{G^0_{36}}{G^0_{37}}} < (v_1)^{(7)} < (\bar{v}_1)^{(7)}$

$$v^{(7)}(t) \ge \frac{(v_1)^{(7)} + (C)^{(7)}(v_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((v_1)^{(7)} - (v_0)^{(7)}\right)t\right]}}{1 + (C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((v_1)^{(7)} - (v_0)^{(7)}\right)t\right]}} \quad , \quad \boxed{(C)^{(7)} = \frac{(v_1)^{(7)} - (v_0)^{(7)}}{(v_0)^{(7)} - (v_2)^{(7)}}}$$

it follows $(v_0)^{(7)} \le v^{(7)}(t) \le (v_1)^{(7)}$

In the same manner , we get

$$v^{(7)}(t) \le \frac{(\bar{v}_1)^{(7)} + (\bar{C})^{(7)}(\bar{v}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}}{1 + (\bar{C})^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(7)} = \frac{(\bar{v}_1)^{(7)} - (v_0)^{(7)}}{(v_0)^{(7)} - (\bar{v}_2)^{(7)}}}$$

From which we deduce $(v_0)^{(7)} \le v^{(7)}(t) \le (\bar{v}_1)^{(7)}$

If $0 < (v_1)^{(7)} < (v_0)^{(7)} = \frac{G^0_{36}}{G^0_{37}} < (\bar{v}_1)^{(7)}$ we find like in the previous case,

$$(v_1)^{(7)} \le \frac{(v_1)^{(7)} + (C)^{(7)}(v_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((v_1)^{(7)} - (v_2)^{(7)}\right)t\right]}}{1 + (C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((v_1)^{(7)} - (v_2)^{(7)}\right)t\right]}} \le v^{(7)}(t) \le$$

$$\frac{(\bar{v}_1)^{(7)} + (\bar{C})^{(7)}(\bar{v}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}}{1 + (C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}} \le (\bar{v}_1)^{(7)}$$

If $0 < (v_1)^{(7)} \le (\bar{v}_1)^{(7)} \le \boxed{(v_0)^{(7)} = \frac{G^0_{36}}{G^0_{37}}}$ , we obtain

$$(v_1)^{(7)} \le v^{(7)}(t) \le \frac{(\bar{v}_1)^{(7)} + (\bar{C})^{(7)}(\bar{v}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}}{1 + (C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}} \le (v_0)^{(7)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $v^{(7)}(t)$ :-

$(m_2)^{(7)} \leq v^{(7)}(t) \leq (m_1)^{(7)}$, $\boxed{v^{(7)}(t) = \frac{G_{36}(t)}{G_{37}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(7)}(t)$ :-

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$(\mu_2)^{(7)} \leq u^{(7)}(t) \leq (\mu_1)^{(7)}$, $\boxed{u^{(7)}(t) = \frac{T_{36}(t)}{T_{37}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{36}^{''})^{(7)} = (a_{37}^{''})^{(7)}, then\ (\sigma_1)^{(7)} = (\sigma_2)^{(7)}$ and in this case $(\nu_1)^{(7)} = (\bar{\nu}_1)^{(7)}$ if in addition $(\nu_0)^{(7)} = (\nu_1)^{(7)}$ then $v^{(7)}(t) = (\nu_0)^{(7)}$ and as a consequence $G_{36}(t) = (\nu_0)^{(7)} G_{37}(t)$ **this also defines** $(\nu_0)^{(7)}$ **for the special case** .

Analogously if $(b_{36}^{''})^{(7)} = (b_{37}^{''})^{(7)}, then\ (\tau_1)^{(7)} = (\tau_2)^{(7)}$ and then $(u_1)^{(7)} = (\bar{u}_1)^{(7)}$if in addition $(u_0)^{(7)} = (u_1)^{(7)}$ then $T_{36}(t) = (u_0)^{(7)} T_{37}(t)$ This is an important consequence of the relation between $(\nu_1)^{(7)}$ and $(\bar{\nu}_1)^{(7)}$, **and definition of** $(u_0)^{(7)}$.

**Proof** : From global equations we obtain

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$$\frac{dv^{(8)}}{dt} = (a_{40})^{(8)} - \left((a_{40}^{'})^{(8)} - (a_{41}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41}, t)\right) - (a_{41}^{''})^{(8)}(T_{41}, t) v^{(8)} - (a_{41})^{(8)} v^{(8)}$$

**Definition of** $v^{(8)}$ :- $\boxed{v^{(8)} = \frac{G_{40}}{G_{41}}}$

It follows

$$-\left((a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_2)^{(8)} v^{(8)} - (a_{40})^{(8)}\right) \leq \frac{dv^{(8)}}{dt} \leq -\left((a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_1)^{(8)} v^{(8)} - (a_{40})^{(8)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(8)}, (\nu_0)^{(8)}$ :-

For $0 < \boxed{(\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}} < (\nu_1)^{(8)} < (\bar{\nu}_1)^{(8)}$

$$v^{(8)}(t) \geq \frac{(\nu_1)^{(8)} + (C)^{(8)}(\nu_2)^{(8)} e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}}{1 + (C)^{(8)} e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}} \quad , \quad \boxed{(C)^{(8)} = \frac{(\nu_1)^{(8)} - (\nu_0)^{(8)}}{(\nu_0)^{(8)} - (\nu_2)^{(8)}}}$$

it follows $(\nu_0)^{(8)} \leq \nu^{(8)}(t) \leq (\nu_1)^{(8)}$

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In the same manner , we get

$$\nu^{(8)}(t) \leq \frac{(\bar{\nu}_1)^{(8)}+(\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}}{1+(\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(8)} = \frac{(\bar{\nu}_1)^{(8)}-(\nu_0)^{(8)}}{(\nu_0)^{(8)}-(\bar{\nu}_2)^{(8)}}}$$

From which we deduce $(\nu_0)^{(8)} \leq \nu^{(8)}(t) \leq (\bar{\nu}_8)^{(8)}$

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If $0 < (\nu_1)^{(8)} < (\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0} < (\bar{\nu}_1)^{(8)}$ we find like in the previous case,

$$(\nu_1)^{(8)} \leq \frac{(\nu_1)^{(8)}+(C)^{(8)}(\nu_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)}-(\nu_2)^{(8)}\right)t\right]}}{1+(C)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)}-(\nu_2)^{(8)}\right)t\right]}} \leq \nu^{(8)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(8)}+(\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}}{1+(\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}} \leq (\bar{\nu}_1)^{(8)}$$

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If $0 < (\nu_1)^{(8)} \leq (\bar{\nu}_1)^{(8)} \leq \boxed{(\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}}$ , we obtain

$$(\nu_1)^{(8)} \leq \nu^{(8)}(t) \leq \frac{(\bar{\nu}_1)^{(8)}+(C)^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}}{1+(C)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}} \leq (\nu_0)^{(8)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(8)}(t)$ :-

$$(m_2)^{(8)} \leq \nu^{(8)}(t) \leq (m_1)^{(8)}, \quad \boxed{\nu^{(8)}(t) = \frac{G_{40}(t)}{G_{41}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(8)}(t)$ :-

$$(\mu_2)^{(8)} \leq u^{(8)}(t) \leq (\mu_1)^{(8)}, \quad \boxed{u^{(8)}(t) = \frac{T_{40}(t)}{T_{41}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{40}^{''})^{(8)} = (a_{41}^{''})^{(8)}, then\ (\sigma_1)^{(8)} = (\sigma_2)^{(8)}$ and in this case $(\nu_1)^{(8)} = (\bar{\nu}_1)^{(8)}$ if in addition $(\nu_0)^{(8)} = (\nu_1)^{(8)}$ then $\nu^{(8)}(t) = (\nu_0)^{(8)}$ and as a consequence $G_{40}(t) = (\nu_0)^{(8)}G_{41}(t)$ **this also defines** $(\nu_0)^{(8)}$ **for the special case** .

Analogously if $(b_{40}^{''})^{(8)} = (b_{41}^{''})^{(8)}, then\ (\tau_1)^{(8)} = (\tau_2)^{(8)}$ and then $(u_1)^{(8)} = (\bar{u}_1)^{(8)}$ if in addition $(u_0)^{(8)} = (u_1)^{(8)}$ then $T_{40}(t) = (u_0)^{(8)}T_{41}(t)$ This is an important

consequence of the relation between $(\nu_1)^{(8)}$ and $(\bar{\nu}_1)^{(8)}$, **and definition of** $(u_0)^{(8)}$.

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**Proof** : From 99,20,44,22,23,44 we obtain

$$\frac{d\nu^{(9)}}{dt} = (a_{44})^{(9)} - \left((a'_{44})^{(9)} - (a'_{45})^{(9)} + (a''_{44})^{(9)}(T_{45},t)\right) - (a''_{45})^{(9)}(T_{45},t)\nu^{(9)} - (a_{45})^{(9)}\nu^{(9)}$$

**Definition of** $\nu^{(9)}$ :- $\boxed{\nu^{(9)} = \frac{G_{44}}{G_{45}}}$

It follows

$$-\left((a_{45})^{(9)}\left(\nu^{(9)}\right)^2 + (\sigma_2)^{(9)}\nu^{(9)} - (a_{44})^{(9)}\right) \leq \frac{d\nu^{(9)}}{dt} \leq -\left((a_{45})^{(9)}\left(\nu^{(9)}\right)^2 + (\sigma_1)^{(9)}\nu^{(9)} - (a_{44})^{(9)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(9)}, (\nu_0)^{(9)}$ :-

For $0 < \boxed{(\nu_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}}} < (\nu_1)^{(9)} < (\bar{\nu}_1)^{(9)}$

$$\nu^{(9)}(t) \geq \frac{(\nu_1)^{(9)} + (C)^{(9)}(\nu_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_0)^{(9)}\right)t\right]}}{1 + (C)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_0)^{(9)}\right)t\right]}} \quad , \quad \boxed{(C)^{(9)} = \frac{(\nu_1)^{(9)} - (\nu_0)^{(9)}}{(\nu_0)^{(9)} - (\nu_2)^{(9)}}}$$

it follows $(\nu_0)^{(9)} \leq \nu^{(9)}(t) \leq (\nu_9)^{(9)}$

In the same manner , we get

$$\nu^{(9)}(t) \leq \frac{(\bar{\nu}_1)^{(9)} + (\bar{C})^{(9)}(\bar{\nu}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}}{1 + (\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(9)} = \frac{(\bar{\nu}_1)^{(9)} - (\nu_0)^{(9)}}{(\nu_0)^{(9)} - (\bar{\nu}_2)^{(9)}}}$$

From which we deduce $(\nu_0)^{(9)} \leq \nu^{(9)}(t) \leq (\bar{\nu}_1)^{(9)}$

If $0 < (\nu_1)^{(9)} < (\nu_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}} < (\bar{\nu}_1)^{(9)}$ we find like in the previous case,

$$(\nu_1)^{(9)} \leq \frac{(\nu_1)^{(9)} + (C)^{(9)}(\nu_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_2)^{(9)}\right)t\right]}}{1 + (C)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_2)^{(9)}\right)t\right]}} \leq \nu^{(9)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(9)} + (\bar{C})^{(9)}(\bar{\nu}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}}{1 + (\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}} \leq (\bar{\nu}_1)^{(9)}$$

If $0 < (\nu_1)^{(9)} \leq (\bar{\nu}_1)^{(9)} \leq \boxed{(\nu_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}}}$ , we obtain

$$(\nu_1)^{(9)} \leq \nu^{(9)}(t) \leq \frac{(\bar{\nu}_1)^{(9)} + (\bar{C})^{(9)}(\bar{\nu}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}}{1 + (\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}} \leq (\nu_0)^{(9)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(9)}(t)$ :-

$(m_2)^{(9)} \leq \nu^{(9)}(t) \leq (m_1)^{(9)}$, $\boxed{\nu^{(9)}(t) = \frac{G_{44}(t)}{G_{45}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(9)}(t)$ :-

$(\mu_2)^{(9)} \leq u^{(9)}(t) \leq (\mu_1)^{(9)}$, $\boxed{u^{(9)}(t) = \frac{T_{44}(t)}{T_{45}(t)}}$

Now, using this result and replacing it in 99, 20,44,22,23, and 44 we get easily the result stated in the theorem.

**Particular case :**

If $(a_{44}^{''})^{(9)} = (a_{45}^{''})^{(9)}, then\ (\sigma_1)^{(9)} = (\sigma_2)^{(9)}$ and in this case $(\nu_1)^{(9)} = (\bar{\nu}_1)^{(9)}$ if in addition $(\nu_0)^{(9)} = (\nu_1)^{(9)}$ then $\nu^{(9)}(t) = (\nu_0)^{(9)}$ and as a consequence $G_{44}(t) = (\nu_0)^{(9)} G_{45}(t)$ **this also defines** $(\nu_0)^{(9)}$ **for the special case** .

Analogously if $(b_{44}^{''})^{(9)} = (b_{45}^{''})^{(9)}, then\ (\tau_1)^{(9)} = (\tau_2)^{(9)}$ and then $(u_1)^{(9)} = (\bar{u}_1)^{(9)}$if in addition $(u_0)^{(9)} = (u_1)^{(9)}$ then $T_{44}(t) = (u_0)^{(9)} T_{45}(t)$ This is an important consequence of the relation between $(\nu_1)^{(9)}$ and $(\bar{\nu}_1)^{(9)}$, **and definition of** $(u_0)^{(9)}$.

We can prove the following 425

**Theorem :** If $(a_i^{''})^{(1)} and\ (b_i^{''})^{(1)}$ are independent on $t$ , and the conditions with the notations

$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} < 0$

$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a_{13})^{(1)}(p_{13})^{(1)} + (a_{14}^{'})^{(1)}(p_{14})^{(1)} + (p_{13})^{(1)}(p_{14})^{(1)} > 0$

$(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} > 0$ ,

$(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} - (b_{13}^{'})^{(1)}(r_{14})^{(1)} - (b_{14}^{'})^{(1)}(r_{14})^{(1)} + (r_{13})^{(1)}(r_{14})^{(1)} < 0$

$with\ (p_{13})^{(1)}, (r_{14})^{(1)}$ as defined by equation are satisfied , then the system

**Theorem :** If $(a_i^{''})^{(2)} and\ (b_i^{''})^{(2)}$ are independent on t , and the conditions with the notations 426

$(a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} < 0$ 427

$(a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a_{16})^{(2)}(p_{16})^{(2)} + (a_{17}^{'})^{(2)}(p_{17})^{(2)} + (p_{16})^{(2)}(p_{17})^{(2)} > 0$ 428

$(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} > 0$ , 429

$(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} - (b_{16}^{'})^{(2)}(r_{17})^{(2)} - (b_{17}^{'})^{(2)}(r_{17})^{(2)} + (r_{16})^{(2)}(r_{17})^{(2)} < 0$ 430

$with\ (p_{16})^{(2)}, (r_{17})^{(2)}$ as defined by equation are satisfied , then the system

**Theorem :** If $(a_i^{''})^{(3)} and\ (b_i^{''})^{(3)}$ are independent on $t$ , and the conditions with the notations 431

$(a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} < 0$

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| $(a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a_{20})^{(3)}(p_{20})^{(3)} + (a_{21}^{'})^{(3)}(p_{21})^{(3)} + (p_{20})^{(3)}(p_{21})^{(3)} > 0$<br>$(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} > 0$ ,<br>$(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} - (b_{20}^{'})^{(3)}(r_{21})^{(3)} - (b_{21}^{'})^{(3)}(r_{21})^{(3)} + (r_{20})^{(3)}(r_{21})^{(3)} < 0$<br>$with\ \ (p_{20})^{(3)}, (r_{21})^{(3)}$ as defined by equation are satisfied , then the system | |
| We can prove the following<br>**Theorem :** If $(a_i^{''})^{(4)} and\ (b_i^{''})^{(4)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} < 0$<br>$(a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a_{24})^{(4)}(p_{24})^{(4)} + (a_{25}^{'})^{(4)}(p_{25})^{(4)} + (p_{24})^{(4)}(p_{25})^{(4)} > 0$<br>$(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} > 0$ ,<br>$(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} - (b_{24}^{'})^{(4)}(r_{25})^{(4)} - (b_{25}^{'})^{(4)}(r_{25})^{(4)} + (r_{24})^{(4)}(r_{25})^{(4)} < 0$<br>$with\ \ (p_{24})^{(4)}, (r_{25})^{(4)}$ as defined by equation are satisfied , then the system | 432 |
| **Theorem :** If $(a_i^{''})^{(5)} and\ (b_i^{''})^{(5)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} < 0$<br>$(a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a_{28})^{(5)}(p_{28})^{(5)} + (a_{29}^{'})^{(5)}(p_{29})^{(5)} + (p_{28})^{(5)}(p_{29})^{(5)} > 0$<br>$(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} > 0$ ,<br>$(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} - (b_{28}^{'})^{(5)}(r_{29})^{(5)} - (b_{29}^{'})^{(5)}(r_{29})^{(5)} + (r_{28})^{(5)}(r_{29})^{(5)} < 0$<br>$with\ \ (p_{28})^{(5)}, (r_{29})^{(5)}$ as defined by equation are satisfied , then the system | 433 |
| **Theorem** If $(a_i^{''})^{(6)} and\ (b_i^{''})^{(6)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} < 0$<br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32})^{(6)}(p_{32})^{(6)} + (a_{33}^{'})^{(6)}(p_{33})^{(6)} + (p_{32})^{(6)}(p_{33})^{(6)} > 0$<br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} > 0$ ,<br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} - (b_{32}^{'})^{(6)}(r_{33})^{(6)} - (b_{33}^{'})^{(6)}(r_{33})^{(6)} + (r_{32})^{(6)}(r_{33})^{(6)} < 0$<br>$with\ \ (p_{32})^{(6)}, (r_{33})^{(6)}$ as defined by equation are satisfied , then the system | 434 |
| **Theorem :** If $(a_i^{''})^{(7)} and\ (b_i^{''})^{(7)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} < 0$ | 435 |

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| $(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36})^{(7)}(p_{36})^{(7)} + (a_{37}^{'})^{(7)}(p_{37})^{(7)} + (p_{36})^{(7)}(p_{37})^{(7)} > 0$<br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} > 0$ ,<br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} - (b_{36}^{'})^{(7)}(r_{37})^{(7)} - (b_{37}^{'})^{(7)}(r_{37})^{(7)} + (r_{36})^{(7)}(r_{37})^{(7)} < 0$<br>$with\ \ (p_{36})^{(7)}, (r_{37})^{(7)}$ as defined by equation are satisfied , then the system | |
| **Theorem :** If $(a_i^{''})^{(8)} and\ (b_i^{''})^{(8)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} < 0$<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40})^{(8)}(p_{40})^{(8)} + (a_{41}^{'})^{(8)}(p_{41})^{(8)} + (p_{40})^{(8)}(p_{41})^{(8)} > 0$<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} > 0$ ,<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} - (b_{40}^{'})^{(8)}(r_{41})^{(8)} - (b_{41}^{'})^{(8)}(r_{41})^{(8)} + (r_{40})^{(8)}(r_{41})^{(8)} < 0$<br>$with\ \ (p_{40})^{(8)}, (r_{41})^{(8)}$ as defined by equation are satisfied , then the system | 436 |
| **Theorem :** If $(a_i^{''})^{(9)} and\ (b_i^{''})^{(9)}$ are independent on $t$ , and the conditions (with the notations 45,46,27,28)<br>$(a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} < 0$<br>$(a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44})^{(9)}(p_{44})^{(9)} + (a_{45}^{'})^{(9)}(p_{45})^{(9)} + (p_{44})^{(9)}(p_{45})^{(9)} > 0$<br>$(b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} > 0$ ,<br>$(b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} - (b_{44}^{'})^{(9)}(r_{45})^{(9)} - (b_{45}^{'})^{(9)}(r_{45})^{(9)} + (r_{44})^{(9)}(r_{45})^{(9)} < 0$<br>$with\ \ (p_{44})^{(9)}, (r_{45})^{(9)}$ as defined by equation 45 are satisfied , then the system | 436 A |
| $(a_{13})^{(1)}G_{14} - \left[(a_{13}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14})\right]G_{13} = 0$ | 437 |
| $(a_{14})^{(1)}G_{13} - \left[(a_{14}^{'})^{(1)} + (a_{14}^{''})^{(1)}(T_{14})\right]G_{14} = 0$ | 438 |
| $(a_{15})^{(1)}G_{14} - \left[(a_{15}^{'})^{(1)} + (a_{15}^{''})^{(1)}(T_{14})\right]G_{15} = 0$ | 439 |
| $(b_{13})^{(1)}T_{14} - [(b_{13}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G)\ ]T_{13} = 0$ | 440 |
| $(b_{14})^{(1)}T_{13} - [(b_{14}^{'})^{(1)} - (b_{14}^{''})^{(1)}(G)\ ]T_{14} = 0$ | 441 |
| $(b_{15})^{(1)}T_{14} - [(b_{15}^{'})^{(1)} - (b_{15}^{''})^{(1)}(G)\ ]T_{15} = 0$ | 442 |
| has a unique positive solution , which is an equilibrium solution for the system | |

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| $(a_{16})^{(2)}G_{17} - \left[(a_{16}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17})\right]G_{16} = 0$ | 443 |
| $(a_{17})^{(2)}G_{16} - \left[(a_{17}^{'})^{(2)} + (a_{17}^{''})^{(2)}(T_{17})\right]G_{17} = 0$ | 444 |
| $(a_{18})^{(2)}G_{17} - \left[(a_{18}^{'})^{(2)} + (a_{18}^{''})^{(2)}(T_{17})\right]G_{18} = 0$ | 445 |
| $(b_{16})^{(2)}T_{17} - [(b_{16}^{'})^{(2)} - (b_{16}^{''})^{(2)}(G_{19})\,]T_{16} = 0$ | 446 |
| $(b_{17})^{(2)}T_{16} - [(b_{17}^{'})^{(2)} - (b_{17}^{''})^{(2)}(G_{19})\,]T_{17} = 0$ | 447 |
| $(b_{18})^{(2)}T_{17} - [(b_{18}^{'})^{(2)} - (b_{18}^{''})^{(2)}(G_{19})\,]T_{18} = 0$ | 448 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{20})^{(3)}G_{21} - \left[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21})\right]G_{20} = 0$ | 449 |
| $(a_{21})^{(3)}G_{20} - \left[(a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)}(T_{21})\right]G_{21} = 0$ | 450 |
| $(a_{22})^{(3)}G_{21} - \left[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21})\right]G_{22} = 0$ | 451 |
| $(b_{20})^{(3)}T_{21} - [(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23})\,]T_{20} = 0$ | 452 |
| $(b_{21})^{(3)}T_{20} - [(b_{21}^{'})^{(3)} - (b_{21}^{''})^{(3)}(G_{23})\,]T_{21} = 0$ | 453 |
| $(b_{22})^{(3)}T_{21} - [(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23})\,]T_{22} = 0$ | 454 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25})\right]G_{24} = 0$ | 455 |
| $(a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25})\right]G_{25} = 0$ | 456 |
| $(a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25})\right]G_{26} = 0$ | 457 |
| $(b_{24})^{(4)}T_{25} - [(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}\left((G_{27})\right)\,]T_{24} = 0$ | 458 |
| $(b_{25})^{(4)}T_{24} - [(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}\left((G_{27})\right)\,]T_{25} = 0$ | 459 |
| $(b_{26})^{(4)}T_{25} - [(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}\left((G_{27})\right)\,]T_{26} = 0$ | 460 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29})\right]G_{28} = 0$ | 461 |
| $(a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29})\right]G_{29} = 0$ | 462 |

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| $(a_{30})^{(5)}G_{29} - \left[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29})\right]G_{30} = 0$ | 463 |
| $(b_{28})^{(5)}T_{29} - [(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}(G_{31})\,]T_{28} = 0$ | 464 |
| $(b_{29})^{(5)}T_{28} - [(b_{29}^{'})^{(5)} - (b_{29}^{''})^{(5)}(G_{31})\,]T_{29} = 0$ | 465 |
| $(b_{30})^{(5)}T_{29} - [(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}(G_{31})\,]T_{30} = 0$ | 466 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{32})^{(6)}G_{33} - \left[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33})\right]G_{32} = 0$ | 467 |
| $(a_{33})^{(6)}G_{32} - \left[(a_{33}^{'})^{(6)} + (a_{33}^{''})^{(6)}(T_{33})\right]G_{33} = 0$ | 468 |
| $(a_{34})^{(6)}G_{33} - \left[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33})\right]G_{34} = 0$ | 469 |
| $(b_{32})^{(6)}T_{33} - [(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}(G_{35})\,]T_{32} = 0$ | 470 |
| $(b_{33})^{(6)}T_{32} - [(b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}(G_{35})\,]T_{33} = 0$ | 471 |
| $(b_{34})^{(6)}T_{33} - [(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}(G_{35})\,]T_{34} = 0$ | 472 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{36})^{(7)}G_{37} - \left[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37})\right]G_{36} = 0$ | 473 |
| $(a_{37})^{(7)}G_{36} - \left[(a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37})\right]G_{37} = 0$ | 474 |
| $(a_{38})^{(7)}G_{37} - \left[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37})\right]G_{38} = 0$ | 475 |
| $(b_{36})^{(7)}T_{37} - [(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}(G_{39})]T_{36} = 0$ | 476 |
| $(b_{37})^{(7)}T_{36} - [(b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}(G_{39})\,]T_{37} = 0$ | 477 |
| $(b_{38})^{(7)}T_{37} - [(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}(G_{39})\,]T_{38} = 0$ | 478 |
| $(a_{40})^{(8)}G_{41} - \left[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41})\right]G_{40} = 0$ | 479 |
| $(a_{41})^{(8)}G_{40} - \left[(a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41})\right]G_{41} = 0$ | 480 |

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| $(a_{42})^{(8)}G_{41} - \left[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41})\right]G_{42} = 0$ | 481 |
| $(b_{40})^{(8)}T_{41} - [(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}(G_{43})\,]T_{40} = 0$ | 482 |
| $(b_{41})^{(8)}T_{40} - [(b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}(G_{43})\,]T_{41} = 0$ | 483 |
| $(b_{42})^{(8)}T_{41} - [(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}(G_{43})\,]T_{42} = 0$ | 484 |
| $(a_{44})^{(9)}G_{45} - \left[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45})\right]G_{44} = 0$ | 484 A |
| $(a_{45})^{(9)}G_{44} - \left[(a_{45}^{'})^{(9)} + (a_{45}^{''})^{(9)}(T_{45})\right]G_{45} = 0$ | |
| $(a_{46})^{(9)}G_{45} - \left[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45})\right]G_{46} = 0$ | |
| $(b_{44})^{(9)}T_{45} - [(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}(G_{47})\,]T_{44} = 0$ | |
| $(b_{45})^{(9)}T_{44} - [(b_{45}^{'})^{(9)} - (b_{45}^{''})^{(9)}(G_{47})\,]T_{45} = 0$ | |
| $(b_{46})^{(9)}T_{45} - [(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}(G_{47})\,]T_{46} = 0$ | |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{13}, G_{14}$ if<br><br>$F(T) = (a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a_{13}^{'})^{(1)}(a_{14}^{''})^{(1)}(T_{14}) + (a_{14}^{'})^{(1)}(a_{13}^{''})^{(1)}(T_{14}) + (a_{13}^{''})^{(1)}(T_{14})(a_{14}^{''})^{(1)}(T_{14}) = 0$ | 485 |
| **Proof:**<br><br>(e) Indeed the first two equations have a nontrivial solution $G_{16}, G_{17}$ if<br><br>$\mathrm{F}(T_{19}) = (a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a_{16}^{'})^{(2)}(a_{17}^{''})^{(2)}(T_{17}) + (a_{17}^{'})^{(2)}(a_{16}^{''})^{(2)}(T_{17}) + (a_{16}^{''})^{(2)}(T_{17})(a_{17}^{''})^{(2)}(T_{17}) = 0$ | 486 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{20}, G_{21}$ if<br><br>$F(T_{23}) = (a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a_{20}^{'})^{(3)}(a_{21}^{''})^{(3)}(T_{21}) + (a_{21}^{'})^{(3)}(a_{20}^{''})^{(3)}(T_{21}) + (a_{20}^{''})^{(3)}(T_{21})(a_{21}^{''})^{(3)}(T_{21}) = 0$ | 487 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{24}, G_{25}$ if<br><br>$F(T_{27}) = (a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a_{24}^{'})^{(4)}(a_{25}^{''})^{(4)}(T_{25}) + (a_{25}^{'})^{(4)}(a_{24}^{''})^{(4)}(T_{25}) +$ | 488 |

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| $(a_{24}^{''})^{(4)}(T_{25})(a_{25}^{''})^{(4)}(T_{25}) = 0$ | |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{28}, G_{29}$ if<br><br>$F(T_{31}) = (a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a_{28}^{'})^{(5)}(a_{29}^{''})^{(5)}(T_{29}) + (a_{29}^{'})^{(5)}(a_{28}^{''})^{(5)}(T_{29}) + (a_{28}^{''})^{(5)}(T_{29})(a_{29}^{''})^{(5)}(T_{29}) = 0$ | 489 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{32}, G_{33}$ if<br><br>$F(T_{35}) = (a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32}^{'})^{(6)}(a_{33}^{''})^{(6)}(T_{33}) + (a_{33}^{'})^{(6)}(a_{32}^{''})^{(6)}(T_{33}) + (a_{32}^{''})^{(6)}(T_{33})(a_{33}^{''})^{(6)}(T_{33}) = 0$ | 490 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{36}, G_{37}$ if<br><br>$F(T_{39}) = (a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36}^{'})^{(7)}(a_{37}^{''})^{(7)}(T_{37}) + (a_{37}^{'})^{(7)}(a_{36}^{''})^{(7)}(T_{37}) + (a_{36}^{''})^{(7)}(T_{37})(a_{37}^{''})^{(7)}(T_{37}) = 0$ | 491 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{40}, G_{41}$ if<br><br>$F(T_{43}) = (a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40}^{'})^{(8)}(a_{41}^{''})^{(8)}(T_{41}) + (a_{41}^{'})^{(8)}(a_{40}^{''})^{(8)}(T_{41}) + (a_{40}^{''})^{(8)}(T_{41})(a_{41}^{''})^{(8)}(T_{41}) = 0$ | 492 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{44}, G_{45}$ if<br><br>$F(T_{47}) = (a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44}^{'})^{(9)}(a_{45}^{''})^{(9)}(T_{45}) + (a_{45}^{'})^{(9)}(a_{44}^{''})^{(9)}(T_{45}) + (a_{44}^{''})^{(9)}(T_{45})(a_{45}^{''})^{(9)}(T_{45}) = 0$ | 492 A |
| **Definition and uniqueness of** $T_{14}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{14})$ being increasing, it follows that there exists a unique $T_{14}^{*}$ for which $f(T_{14}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{13} = \frac{(a_{13})^{(1)}G_{14}}{[(a_{13}^{'})^{(1)}+(a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15} = \frac{(a_{15})^{(1)}G_{14}}{[(a_{15}^{'})^{(1)}+(a_{15}^{''})^{(1)}(T_{14}^{*})]}$ | 493 |
| **Definition and uniqueness of** $T_{17}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(2)}(T_{17})$ being increasing, it follows that there exists a unique $T_{17}^{*}$ for which $f(T_{17}^{*}) = 0$. With this value , we obtain from the three first | 494 |

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| equations | |
| $G_{16} = \frac{(a_{16})^{(2)}G_{17}}{[(a_{16}^{'})^{(2)}+(a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18} = \frac{(a_{18})^{(2)}G_{17}}{[(a_{18}^{'})^{(2)}+(a_{18}^{''})^{(2)}(T_{17}^{*})]}$ | 495 |
| **Definition and uniqueness of** $T_{21}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{21})$ being increasing, it follows that there exists a unique $T_{21}^{*}$ for which $f(T_{21}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{20} = \frac{(a_{20})^{(3)}G_{21}}{[(a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21}^{*})]}$ , $G_{22} = \frac{(a_{22})^{(3)}G_{21}}{[(a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21}^{*})]}$ | 496 |
| **Definition and uniqueness of** $T_{25}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(4)}(T_{25})$ being increasing, it follows that there exists a unique $T_{25}^{*}$ for which $f(T_{25}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{24} = \frac{(a_{24})^{(4)}G_{25}}{[(a_{24}^{'})^{(4)}+(a_{24}^{''})^{(4)}(T_{25}^{*})]}$ , $G_{26} = \frac{(a_{26})^{(4)}G_{25}}{[(a_{26}^{'})^{(4)}+(a_{26}^{''})^{(4)}(T_{25}^{*})]}$ | 497 |
| **Definition and uniqueness of** $T_{29}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(5)}(T_{29})$ being increasing, it follows that there exists a unique $T_{29}^{*}$ for which $f(T_{29}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{28} = \frac{(a_{28})^{(5)}G_{29}}{[(a_{28}^{'})^{(5)}+(a_{28}^{''})^{(5)}(T_{29}^{*})]}$ , $G_{30} = \frac{(a_{30})^{(5)}G_{29}}{[(a_{30}^{'})^{(5)}+(a_{30}^{''})^{(5)}(T_{29}^{*})]}$ | 498 |
| **Definition and uniqueness of** $T_{33}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(6)}(T_{33})$ being increasing, it follows that there exists a unique $T_{33}^{*}$ for which $f(T_{33}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{32} = \frac{(a_{32})^{(6)}G_{33}}{[(a_{32}^{'})^{(6)}+(a_{32}^{''})^{(6)}(T_{33}^{*})]}$ , $G_{34} = \frac{(a_{34})^{(6)}G_{33}}{[(a_{34}^{'})^{(6)}+(a_{34}^{''})^{(6)}(T_{33}^{*})]}$ | 499 |
| **Definition and uniqueness of** $T_{37}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(7)}(T_{37})$ being increasing, it follows that there exists a unique $T_{37}^{*}$ for which $f(T_{37}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{36} = \frac{(a_{36})^{(7)}G_{37}}{[(a_{36}^{'})^{(7)}+(a_{36}^{''})^{(7)}(T_{37}^{*})]}$ , $G_{38} = \frac{(a_{38})^{(7)}G_{37}}{[(a_{38}^{'})^{(7)}+(a_{38}^{''})^{(7)}(T_{37}^{*})]}$ | 500 |

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| **Definition and uniqueness of** $T_{41}^{*}$ :-<br><br>After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(8)}(T_{41})$ being increasing, it follows that there exists a unique $T_{41}^{*}$ for which $f(T_{41}^{*})=0$. With this value , we obtain from the three first equations<br><br>$G_{40}=\frac{(a_{40})^{(8)}G_{41}}{[(a_{40}^{'})^{(8)}+(a_{40}^{''})^{(8)}(T_{41}^{*})]}$ , $G_{42}=\frac{(a_{42})^{(8)}G_{41}}{[(a_{42}^{'})^{(8)}+(a_{42}^{''})^{(8)}(T_{41}^{*})]}$ | 501 |
| **Definition and uniqueness of** $T_{45}^{*}$ :-<br><br>After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(9)}(T_{45})$ being increasing, it follows that there exists a unique $T_{45}^{*}$ for which $f(T_{45}^{*})=0$. With this value , we obtain from the three first equations<br><br>$G_{44}=\frac{(a_{44})^{(9)}G_{45}}{[(a_{44}^{'})^{(9)}+(a_{44}^{''})^{(9)}(T_{45}^{*})]}$ , $G_{46}=\frac{(a_{46})^{(9)}G_{45}}{[(a_{46}^{'})^{(9)}+(a_{46}^{''})^{(9)}(T_{45}^{*})]}$ | 501 A |
| By the same argument, the equations admit solutions $G_{13}, G_{14}$ if<br>$\varphi(G)=(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)}-(b_{13})^{(1)}(b_{14})^{(1)}-$<br>$[(b_{13}^{'})^{(1)}(b_{14}^{''})^{(1)}(G)+(b_{14}^{'})^{(1)}(b_{13}^{''})^{(1)}(G)]+(b_{13}^{''})^{(1)}(G)(b_{14}^{''})^{(1)}(G)=0$<br><br>Where in $G(G_{13},G_{14},G_{15}), G_{13}, G_{15}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{14}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{14}^{*}$ such that $\varphi(G^{*})=0$ | 502 |
| By the same argument, the equations admit solutions $G_{16}, G_{17}$ if<br><br>$\varphi(G_{19})=(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)}-(b_{16})^{(2)}(b_{17})^{(2)}-$<br>$[(b_{16}^{'})^{(2)}(b_{17}^{''})^{(2)}(G_{19})+(b_{17}^{'})^{(2)}(b_{16}^{''})^{(2)}(G_{19})]+(b_{16}^{''})^{(2)}(G_{19})(b_{17}^{''})^{(2)}(G_{19})=0$ | 503 |
| Where in $(G_{19})(G_{16},G_{17},G_{18}), G_{16}, G_{18}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{17}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{14}^{*}$ such that $\varphi((G_{19})^{*})=0$ | 504 |
| By the same argument, the equations admit solutions $G_{20}, G_{21}$ if<br><br>$\varphi(G_{23})=(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)}-(b_{20})^{(3)}(b_{21})^{(3)}-$<br>$[(b_{20}^{'})^{(3)}(b_{21}^{''})^{(3)}(G_{23})+(b_{21}^{'})^{(3)}(b_{20}^{''})^{(3)}(G_{23})]+(b_{20}^{''})^{(3)}(G_{23})(b_{21}^{''})^{(3)}(G_{23})=0$<br><br>Where in $G_{23}(G_{20},G_{21},G_{22}), G_{20}, G_{22}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{21}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{21}^{*}$ such that $\varphi((G_{23})^{*})=0$ | 505 |

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| By the same argument, the equations admit solutions $G_{24}, G_{25}$ if<br><br>$\varphi(G_{27}) = (b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} -$<br><br>$\left[(b_{24}^{'})^{(4)}(b_{25}^{''})^{(4)}(G_{27}) + (b_{25}^{'})^{(4)}(b_{24}^{''})^{(4)}(G_{27})\right] + (b_{24}^{''})^{(4)}(G_{27})(b_{25}^{''})^{(4)}(G_{27}) = 0$<br><br>Where in $(G_{27})(G_{24}, G_{25}, G_{26}), G_{24}, G_{26}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{25}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{25}^{*}$ such that $\varphi((G_{27})^{*}) = 0$ | 506 |
| By the same argument, the equations admit solutions $G_{28}, G_{29}$ if<br>$\varphi(G_{31}) = (b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} -$<br><br>$\left[(b_{28}^{'})^{(5)}(b_{29}^{''})^{(5)}(G_{31}) + (b_{29}^{'})^{(5)}(b_{28}^{''})^{(5)}(G_{31})\right] + (b_{28}^{''})^{(5)}(G_{31})(b_{29}^{''})^{(5)}(G_{31}) = 0$<br><br>Where in $(G_{31})(G_{28}, G_{29}, G_{30}), G_{28}, G_{30}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{29}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{29}^{*}$ such that $\varphi((G_{31})^{*}) = 0$ | 507 |
| By the same argument, the equations admit solutions $G_{32}, G_{33}$ if<br>$\varphi(G_{35}) = (b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} -$<br><br>$\left[(b_{32}^{'})^{(6)}(b_{33}^{''})^{(6)}(G_{35}) + (b_{33}^{'})^{(6)}(b_{32}^{''})^{(6)}(G_{35})\right] + (b_{32}^{''})^{(6)}(G_{35})(b_{33}^{''})^{(6)}(G_{35}) = 0$<br><br>Where in $(G_{35})(G_{32}, G_{33}, G_{34}), G_{32}, G_{34}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{33}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{33}^{*}$ such that $\varphi({G_{35}}^{*}) = 0$ | 508 |
| By the same argument, the equations admit solutions $G_{36}, G_{37}$ if<br>$\varphi(G_{39}) = (b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} -$<br>$\left[(b_{36}^{'})^{(7)}(b_{37}^{''})^{(7)}(G_{39}) + (b_{37}^{'})^{(7)}(b_{36}^{''})^{(7)}(G_{39})\right] + (b_{36}^{''})^{(7)}(G_{39})(b_{37}^{''})^{(7)}(G_{39}) = 0$<br><br>Where in $(G_{39})(G_{36}, G_{37}, G_{38}), G_{36}, G_{38}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{37}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{37}^{*}$ such that $\varphi({G_{39}}^{*}) = 0$ | 509 |
| By the same argument, the equations admit solutions $G_{40}, G_{41}$ if<br>$\varphi(G_{43}) = (b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} -$<br>$\left[(b_{40}^{'})^{(8)}(b_{41}^{''})^{(8)}(G_{43}) + (b_{41}^{'})^{(8)}(b_{40}^{''})^{(8)}(G_{43})\right] + (b_{40}^{''})^{(8)}(G_{43})(b_{41}^{''})^{(8)}(G_{43}) = 0$<br><br>Where in $(G_{43})(G_{40}, G_{41}, G_{42}), G_{40}, G_{42}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{41}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{41}^{*}$ such that $\varphi({G_{43}}^{*}) = 0$ | 510 |
| By the same argument, the equations 92,93 admit solutions $G_{44}, G_{45}$ if<br>$\varphi(G_{47}) = (b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} -$ | |

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| $\left[(b_{44}^{'})^{(9)}(b_{45}^{''})^{(9)}(G_{47}) + (b_{45}^{'})^{(9)}(b_{44}^{''})^{(9)}(G_{47})\right]+(b_{44}^{''})^{(9)}(G_{47})(b_{45}^{''})^{(9)}(G_{47}) = 0$<br><br>Where in $(G_{47})(G_{44}, G_{45}, G_{46}), G_{44}, G_{46}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{45}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{45}^{*}$ such that $\varphi((G_{47})^{*}) = 0$ | |
| Finally we obtain the unique solution<br><br>$G_{14}^{*}$ given by $\varphi(G^{*}) = 0$ , $T_{14}^{*}$ given by $f(T_{14}^{*}) = 0$ and<br><br>$G_{13}^{*} = \frac{(a_{13})^{(1)}G_{14}^{*}}{[(a_{13}^{'})^{(1)}+(a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15}^{*} = \frac{(a_{15})^{(1)}G_{14}^{*}}{[(a_{15}^{'})^{(1)}+(a_{15}^{''})^{(1)}(T_{14}^{*})]}$<br><br>$T_{13}^{*} = \frac{(b_{13})^{(1)}T_{14}^{*}}{[(b_{13}^{'})^{(1)}-(b_{13}^{''})^{(1)}(G^{*})]}$ , $T_{15}^{*} = \frac{(b_{15})^{(1)}T_{14}^{*}}{[(b_{15}^{'})^{(1)}-(b_{15}^{''})^{(1)}(G^{*})]}$<br><br>Obviously, these values represent an equilibrium solution | 511 |
| Finally we obtain the unique solution | |
| $G_{17}^{*}$ given by $\varphi((G_{19})^{*}) = 0$ , $T_{17}^{*}$ given by $f(T_{17}^{*}) = 0$ and | 512 |
| $G_{16}^{*} = \frac{(a_{16})^{(2)}G_{17}^{*}}{[(a_{16}^{'})^{(2)}+(a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18}^{*} = \frac{(a_{18})^{(2)}G_{17}^{*}}{[(a_{18}^{'})^{(2)}+(a_{18}^{''})^{(2)}(T_{17}^{*})]}$ | 513 |
| $T_{16}^{*} = \frac{(b_{16})^{(2)}T_{17}^{*}}{[(b_{16}^{'})^{(2)}-(b_{16}^{''})^{(2)}((G_{19})^{*})]}$ , $T_{18}^{*} = \frac{(b_{18})^{(2)}T_{17}^{*}}{[(b_{18}^{'})^{(2)}-(b_{18}^{''})^{(2)}((G_{19})^{*})]}$ | 514 |
| Obviously, these values represent an equilibrium solution | |
| Finally we obtain the unique solution<br><br>$G_{21}^{*}$ given by $\varphi((G_{23})^{*}) = 0$ , $T_{21}^{*}$ given by $f(T_{21}^{*}) = 0$ and<br><br>$G_{20}^{*} = \frac{(a_{20})^{(3)}G_{21}^{*}}{[(a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21}^{*})]}$ , $G_{22}^{*} = \frac{(a_{22})^{(3)}G_{21}^{*}}{[(a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21}^{*})]}$<br><br>$T_{20}^{*} = \frac{(b_{20})^{(3)}T_{21}^{*}}{[(b_{20}^{'})^{(3)}-(b_{20}^{''})^{(3)}(G_{23}{}^{*})]}$ , $T_{22}^{*} = \frac{(b_{22})^{(3)}T_{21}^{*}}{[(b_{22}^{'})^{(3)}-(b_{22}^{''})^{(3)}(G_{23}{}^{*})]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 515 |
| Finally we obtain the unique solution<br><br>$G_{25}^{*}$ given by $\varphi(G_{27}) = 0$ , $T_{25}^{*}$ given by $f(T_{25}^{*}) = 0$ and<br><br>$G_{24}^{*} = \frac{(a_{24})^{(4)}G_{25}^{*}}{[(a_{24}^{'})^{(4)}+(a_{24}^{''})^{(4)}(T_{25}^{*})]}$ , $G_{26}^{*} = \frac{(a_{26})^{(4)}G_{25}^{*}}{[(a_{26}^{'})^{(4)}+(a_{26}^{''})^{(4)}(T_{25}^{*})]}$ | 516 |
| $T_{24}^{*} = \frac{(b_{24})^{(4)}T_{25}^{*}}{[(b_{24}^{'})^{(4)}-(b_{24}^{''})^{(4)}((G_{27})^{*})]}$ , $T_{26}^{*} = \frac{(b_{26})^{(4)}T_{25}^{*}}{[(b_{26}^{'})^{(4)}-(b_{26}^{''})^{(4)}((G_{27})^{*})]}$ | 517 |

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|---|---|
| Obviously, these values represent an equilibrium solution of global equations | |
| Finally we obtain the unique solution<br><br>$G_{29}^{*}$ given by $\varphi((G_{31})^{*}) = 0$ , $T_{29}^{*}$ given by $f(T_{29}^{*}) = 0$ and<br><br>$G_{28}^{*} = \frac{(a_{28})^{(5)} G_{29}^{*}}{[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}^{*})]}$ , $G_{30}^{*} = \frac{(a_{30})^{(5)} G_{29}^{*}}{[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29}^{*})]}$ | 518 |
| $T_{28}^{*} = \frac{(b_{28})^{(5)} T_{29}^{*}}{[(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}((G_{31})^{*})]}$ , $T_{30}^{*} = \frac{(b_{30})^{(5)} T_{29}^{*}}{[(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}((G_{31})^{*})]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 519 |
| Finally we obtain the unique solution<br><br>$G_{33}^{*}$ given by $\varphi((G_{35})^{*}) = 0$ , $T_{33}^{*}$ given by $f(T_{33}^{*}) = 0$ and<br><br>$G_{32}^{*} = \frac{(a_{32})^{(6)} G_{33}^{*}}{[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33}^{*})]}$ , $G_{34}^{*} = \frac{(a_{34})^{(6)} G_{33}^{*}}{[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33}^{*})]}$ | 520 |
| $T_{32}^{*} = \frac{(b_{32})^{(6)} T_{33}^{*}}{[(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}((G_{35})^{*})]}$ , $T_{34}^{*} = \frac{(b_{34})^{(6)} T_{33}^{*}}{[(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}((G_{35})^{*})]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 521 |
| Finally we obtain the unique solution<br><br>$G_{37}^{*}$ given by $\varphi((G_{39})^{*}) = 0$ , $T_{37}^{*}$ given by $f(T_{37}^{*}) = 0$ and<br><br>$G_{36}^{*} = \frac{(a_{36})^{(7)} G_{37}^{*}}{[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37}^{*})]}$ , $G_{38}^{*} = \frac{(a_{38})^{(7)} G_{37}^{*}}{[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37}^{*})]}$<br><br>$T_{36}^{*} = \frac{(b_{36})^{(7)} T_{37}^{*}}{[(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}((G_{39})^{*})]}$ , $T_{38}^{*} = \frac{(b_{38})^{(7)} T_{37}^{*}}{[(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}((G_{39})^{*})]}$ | 522 |
| Finally we obtain the unique solution<br><br>$G_{41}^{*}$ given by $\varphi((G_{43})^{*}) = 0$ , $T_{41}^{*}$ given by $f(T_{41}^{*}) = 0$ and<br><br>$G_{40}^{*} = \frac{(a_{40})^{(8)} G_{41}^{*}}{[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41}^{*})]}$ , $G_{42}^{*} = \frac{(a_{42})^{(8)} G_{41}^{*}}{[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41}^{*})]}$<br><br>$T_{40}^{*} = \frac{(b_{40})^{(8)} T_{41}^{*}}{[(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}((G_{43})^{*})]}$ , $T_{42}^{*} = \frac{(b_{42})^{(8)} T_{41}^{*}}{[(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}((G_{43})^{*})]}$ | 523 |
| Finally we obtain the unique solution of 89 to 99 | 523 A |

$G_{45}^{*}$ given by $\varphi((G_{47})^{*}) = 0$ , $T_{45}^{*}$ given by $f(T_{45}^{*}) = 0$ and

$$G_{44}^{*} = \frac{(a_{44})^{(9)} G_{45}^{*}}{[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45}^{*})]} \quad , \quad G_{46}^{*} = \frac{(a_{46})^{(9)} G_{45}^{*}}{[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45}^{*})]}$$

$$T_{44}^{*} = \frac{(b_{44})^{(9)} T_{45}^{*}}{[(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}((G_{47})^{*})]} \quad , \quad T_{46}^{*} = \frac{(b_{46})^{(9)} T_{45}^{*}}{[(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}((G_{47})^{*})]}$$

**ASYMPTOTIC STABILITY ANALYSIS** 524

**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(1)}$ *and* $(b_i^{''})^{(1)}$ Belong to $C^{(1)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.

**Proof:** Denote

**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-

$$G_i = G_i^{*} + \mathbb{G}_i \quad , T_i = T_i^{*} + \mathbb{T}_i$$

$$\frac{\partial (a_{14}^{''})^{(1)}}{\partial T_{14}}(T_{14}^{*}) = (q_{14})^{(1)} \quad , \frac{\partial (b_i^{''})^{(1)}}{\partial G_j}(G^{*}) = s_{ij}$$

Then taking into account equations and neglecting the terms of power 2, we obtain

$$\frac{d\mathbb{G}_{13}}{dt} = -\left((a_{13}^{'})^{(1)} + (p_{13})^{(1)}\right)\mathbb{G}_{13} + (a_{13})^{(1)}\mathbb{G}_{14} - (q_{13})^{(1)} G_{13}^{*}\mathbb{T}_{14} \quad 525$$

$$\frac{d\mathbb{G}_{14}}{dt} = -\left((a_{14}^{'})^{(1)} + (p_{14})^{(1)}\right)\mathbb{G}_{14} + (a_{14})^{(1)}\mathbb{G}_{13} - (q_{14})^{(1)} G_{14}^{*}\mathbb{T}_{14} \quad 526$$

$$\frac{d\mathbb{G}_{15}}{dt} = -\left((a_{15}^{'})^{(1)} + (p_{15})^{(1)}\right)\mathbb{G}_{15} + (a_{15})^{(1)}\mathbb{G}_{14} - (q_{15})^{(1)} G_{15}^{*}\mathbb{T}_{14} \quad 527$$

$$\frac{d\mathbb{T}_{13}}{dt} = -\left((b_{13}^{'})^{(1)} - (r_{13})^{(1)}\right)\mathbb{T}_{13} + (b_{13})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(13)(j)} T_{13}^{*}\mathbb{G}_j\right) \quad 528$$

$$\frac{d\mathbb{T}_{14}}{dt} = -\left((b_{14}^{'})^{(1)} - (r_{14})^{(1)}\right)\mathbb{T}_{14} + (b_{14})^{(1)}\mathbb{T}_{13} + \sum_{j=13}^{15}\left(s_{(14)(j)} T_{14}^{*}\mathbb{G}_j\right) \quad 529$$

$$\frac{d\mathbb{T}_{15}}{dt} = -\left((b_{15}^{'})^{(1)} - (r_{15})^{(1)}\right)\mathbb{T}_{15} + (b_{15})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(15)(j)} T_{15}^{*}\mathbb{G}_j\right) \quad 530$$

**ASYMPTOTIC STABILITY ANALYSIS** 531

**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}$ Belong to $C^{(2)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable

Proof: Denote

**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-

$$G_i = G_i^{*} + \mathbb{G}_i \quad , T_i = T_i^{*} + \mathbb{T}_i \quad 532$$

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| $\frac{\partial(a_{17}^{''})^{(2)}}{\partial \mathrm{T}_{17}}(\mathrm{T}_{17}^{*}) = (q_{17})^{(2)}$ , $\frac{\partial(b_i^{''})^{(2)}}{\partial \mathrm{G}_j}((G_{19})^{*}) = s_{ij}$ | 533 |
| taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{\mathrm{d}\mathbb{G}_{16}}{\mathrm{dt}} = -\left((a_{16}^{'})^{(2)} + (p_{16})^{(2)}\right)\mathbb{G}_{16} + (a_{16})^{(2)}\mathbb{G}_{17} - (q_{16})^{(2)}\mathrm{G}_{16}^{*}\mathbb{T}_{17}$ | 534 |
| $\frac{\mathrm{d}\mathbb{G}_{17}}{\mathrm{dt}} = -\left((a_{17}^{'})^{(2)} + (p_{17})^{(2)}\right)\mathbb{G}_{17} + (a_{17})^{(2)}\mathbb{G}_{16} - (q_{17})^{(2)}\mathrm{G}_{17}^{*}\mathbb{T}_{17}$ | 535 |
| $\frac{\mathrm{d}\mathbb{G}_{18}}{\mathrm{dt}} = -\left((a_{18}^{'})^{(2)} + (p_{18})^{(2)}\right)\mathbb{G}_{18} + (a_{18})^{(2)}\mathbb{G}_{17} - (q_{18})^{(2)}\mathrm{G}_{18}^{*}\mathbb{T}_{17}$ | 536 |
| $\frac{\mathrm{d}\mathbb{T}_{16}}{\mathrm{dt}} = -\left((b_{16}^{'})^{(2)} - (r_{16})^{(2)}\right)\mathbb{T}_{16} + (b_{16})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(16)(j)}\mathrm{T}_{16}^{*}\mathbb{G}_j\right)$ | 537 |
| $\frac{\mathrm{d}\mathbb{T}_{17}}{\mathrm{dt}} = -\left((b_{17}^{'})^{(2)} - (r_{17})^{(2)}\right)\mathbb{T}_{17} + (b_{17})^{(2)}\mathbb{T}_{16} + \sum_{j=16}^{18}\left(s_{(17)(j)}\mathrm{T}_{17}^{*}\mathbb{G}_j\right)$ | 538 |
| $\frac{\mathrm{d}\mathbb{T}_{18}}{\mathrm{dt}} = -\left((b_{18}^{'})^{(2)} - (r_{18})^{(2)}\right)\mathbb{T}_{18} + (b_{18})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(18)(j)}\mathrm{T}_{18}^{*}\mathbb{G}_j\right)$ | 539 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(3)}$ *and* $(b_i^{''})^{(3)}$ Belong to $C^{(3)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^{*} + \mathbb{G}_i \quad , T_i = T_i^{*} + \mathbb{T}_i$<br><br>$\frac{\partial(a_{21}^{''})^{(3)}}{\partial T_{21}}(T_{21}^{*}) = (q_{21})^{(3)}$ , $\frac{\partial(b_i^{''})^{(3)}}{\partial G_j}((G_{23})^{*}) = s_{ij}$ | 540 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{20}}{dt} = -\left((a_{20}^{'})^{(3)} + (p_{20})^{(3)}\right)\mathbb{G}_{20} + (a_{20})^{(3)}\mathbb{G}_{21} - (q_{20})^{(3)}G_{20}^{*}\mathbb{T}_{21}$ | 541 |
| $\frac{d\mathbb{G}_{21}}{dt} = -\left((a_{21}^{'})^{(3)} + (p_{21})^{(3)}\right)\mathbb{G}_{21} + (a_{21})^{(3)}\mathbb{G}_{20} - (q_{21})^{(3)}G_{21}^{*}\mathbb{T}_{21}$ | 542 |
| $\frac{d\mathbb{G}_{22}}{dt} = -\left((a_{22}^{'})^{(3)} + (p_{22})^{(3)}\right)\mathbb{G}_{22} + (a_{22})^{(3)}\mathbb{G}_{21} - (q_{22})^{(3)}G_{22}^{*}\mathbb{T}_{21}$ | 543 |
| $\frac{d\mathbb{T}_{20}}{dt} = -\left((b_{20}^{'})^{(3)} - (r_{20})^{(3)}\right)\mathbb{T}_{20} + (b_{20})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(20)(j)}T_{20}^{*}\mathbb{G}_j\right)$ | 544 |
| $\frac{d\mathbb{T}_{21}}{dt} = -\left((b_{21}^{'})^{(3)} - (r_{21})^{(3)}\right)\mathbb{T}_{21} + (b_{21})^{(3)}\mathbb{T}_{20} + \sum_{j=20}^{22}\left(s_{(21)(j)}T_{21}^{*}\mathbb{G}_j\right)$ | 545 |
| $\frac{d\mathbb{T}_{22}}{dt} = -\left((b_{22}^{'})^{(3)} - (r_{22})^{(3)}\right)\mathbb{T}_{22} + (b_{22})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(22)(j)}T_{22}^{*}\mathbb{G}_j\right)$ | 546 |
| **ASYMPTOTIC STABILITY ANALYSIS** | 547 |

**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(4)}$ $and$ $(b_i^{''})^{(4)}$ Belong to $C^{(4)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.

Proof: Denote

**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-

$$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$$

$$\frac{\partial (a_{25}^{''})^{(4)}}{\partial T_{25}}(T_{25}^*) = (q_{25})^{(4)} \quad , \frac{\partial (b_i^{''})^{(4)}}{\partial G_j}((G_{27})^* ) = s_{ij}$$

Then taking into account equations and neglecting the terms of power 2, we obtain

$$\frac{d\mathbb{G}_{24}}{dt} = -\left((a_{24}^{'})^{(4)} + (p_{24})^{(4)}\right)\mathbb{G}_{24} + (a_{24})^{(4)}\mathbb{G}_{25} - (q_{24})^{(4)}G_{24}^*\mathbb{T}_{25} \quad (549)$$

$$\frac{d\mathbb{G}_{25}}{dt} = -\left((a_{25}^{'})^{(4)} + (p_{25})^{(4)}\right)\mathbb{G}_{25} + (a_{25})^{(4)}\mathbb{G}_{24} - (q_{25})^{(4)}G_{25}^*\mathbb{T}_{25} \quad (550)$$

$$\frac{d\mathbb{G}_{26}}{dt} = -\left((a_{26}^{'})^{(4)} + (p_{26})^{(4)}\right)\mathbb{G}_{26} + (a_{26})^{(4)}\mathbb{G}_{25} - (q_{26})^{(4)}G_{26}^*\mathbb{T}_{25} \quad (551)$$

$$\frac{d\mathbb{T}_{24}}{dt} = -\left((b_{24}^{'})^{(4)} - (r_{24})^{(4)}\right)\mathbb{T}_{24} + (b_{24})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(24)(j)}T_{24}^*\mathbb{G}_j\right) \quad (552)$$

$$\frac{d\mathbb{T}_{25}}{dt} = -\left((b_{25}^{'})^{(4)} - (r_{25})^{(4)}\right)\mathbb{T}_{25} + (b_{25})^{(4)}\mathbb{T}_{24} + \sum_{j=24}^{26}\left(s_{(25)(j)}T_{25}^*\mathbb{G}_j\right) \quad (553)$$

$$\frac{d\mathbb{T}_{26}}{dt} = -\left((b_{26}^{'})^{(4)} - (r_{26})^{(4)}\right)\mathbb{T}_{26} + (b_{26})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(26)(j)}T_{26}^*\mathbb{G}_j\right) \quad (554)$$

**ASYMPTOTIC STABILITY ANALYSIS** (555)

**Theorem 5:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(5)}$ $and$ $(b_i^{''})^{(5)}$ Belong to $C^{(5)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.

Proof: Denote

**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- (556)

$$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$$

$$\frac{\partial (a_{29}^{''})^{(5)}}{\partial T_{29}}(T_{29}^*) = (q_{29})^{(5)} \quad , \frac{\partial (b_i^{''})^{(5)}}{\partial G_j}(\,(G_{31})^*\,) = s_{ij}$$

Then taking into account equations and neglecting the terms of power 2, we obtain

$$\frac{d\mathbb{G}_{28}}{dt} = -\left((a_{28}^{'})^{(5)} + (p_{28})^{(5)}\right)\mathbb{G}_{28} + (a_{28})^{(5)}\mathbb{G}_{29} - (q_{28})^{(5)}G_{28}^*\mathbb{T}_{29} \quad (557)$$

$$\frac{d\mathbb{G}_{29}}{dt} = -\left((a_{29}^{'})^{(5)} + (p_{29})^{(5)}\right)\mathbb{G}_{29} + (a_{29})^{(5)}\mathbb{G}_{28} - (q_{29})^{(5)}G_{29}^*\mathbb{T}_{29} \quad (558)$$

$$\frac{d\mathbb{G}_{30}}{dt} = -\left((a_{30}^{'})^{(5)} + (p_{30})^{(5)}\right)\mathbb{G}_{30} + (a_{30})^{(5)}\mathbb{G}_{29} - (q_{30})^{(5)}G_{30}^*\mathbb{T}_{29} \quad (559)$$

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| $\frac{d\mathbb{T}_{28}}{dt} = -\left((b_{28}^{'})^{(5)} - (r_{28})^{(5)}\right)\mathbb{T}_{28} + (b_{28})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(28)(j)}T_{28}^{*}\mathbb{G}_{j}\right)$ | 560 |
| $\frac{d\mathbb{T}_{29}}{dt} = -\left((b_{29}^{'})^{(5)} - (r_{29})^{(5)}\right)\mathbb{T}_{29} + (b_{29})^{(5)}\mathbb{T}_{28} + \sum_{j=28}^{30}\left(s_{(29)(j)}T_{29}^{*}\mathbb{G}_{j}\right)$ | 561 |
| $\frac{d\mathbb{T}_{30}}{dt} = -\left((b_{30}^{'})^{(5)} - (r_{30})^{(5)}\right)\mathbb{T}_{30} + (b_{30})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(30)(j)}T_{30}^{*}\mathbb{G}_{j}\right)$ | 562 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 6:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(6)}$ $and$ $(b_i^{''})^{(6)}$ Belong to $C^{(6)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 563 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial(a_{33}^{''})^{(6)}}{\partial T_{33}}(T_{33}^{*}) = (q_{33})^{(6)}$ , $\frac{\partial(b_i^{''})^{(6)}}{\partial G_j}((G_{35})^{*}) = s_{ij}$ | 564 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{32}}{dt} = -\left((a_{32}^{'})^{(6)} + (p_{32})^{(6)}\right)\mathbb{G}_{32} + (a_{32})^{(6)}\mathbb{G}_{33} - (q_{32})^{(6)}G_{32}^{*}\mathbb{T}_{33}$ | 565 |
| $\frac{d\mathbb{G}_{33}}{dt} = -\left((a_{33}^{'})^{(6)} + (p_{33})^{(6)}\right)\mathbb{G}_{33} + (a_{33})^{(6)}\mathbb{G}_{32} - (q_{33})^{(6)}G_{33}^{*}\mathbb{T}_{33}$ | 566 |
| $\frac{d\mathbb{G}_{34}}{dt} = -\left((a_{34}^{'})^{(6)} + (p_{34})^{(6)}\right)\mathbb{G}_{34} + (a_{34})^{(6)}\mathbb{G}_{33} - (q_{34})^{(6)}G_{34}^{*}\mathbb{T}_{33}$ | 567 |
| $\frac{d\mathbb{T}_{32}}{dt} = -\left((b_{32}^{'})^{(6)} - (r_{32})^{(6)}\right)\mathbb{T}_{32} + (b_{32})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(32)(j)}T_{32}^{*}\mathbb{G}_{j}\right)$ | 568 |
| $\frac{d\mathbb{T}_{33}}{dt} = -\left((b_{33}^{'})^{(6)} - (r_{33})^{(6)}\right)\mathbb{T}_{33} + (b_{33})^{(6)}\mathbb{T}_{32} + \sum_{j=32}^{34}\left(s_{(33)(j)}T_{33}^{*}\mathbb{G}_{j}\right)$ | 569 |
| $\frac{d\mathbb{T}_{34}}{dt} = -\left((b_{34}^{'})^{(6)} - (r_{34})^{(6)}\right)\mathbb{T}_{34} + (b_{34})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(34)(j)}T_{34}^{*}\mathbb{G}_{j}\right)$ | 570 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 7:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(7)}$ $and$ $(b_i^{''})^{(7)}$ Belong to $C^{(7)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 571 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$ | 572 |

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| $\frac{\partial (a_{37}^{''})^{(7)}}{\partial T_{37}}(T_{37}^{*}) = (q_{37})^{(7)}$ , $\frac{\partial (b_i^{''})^{(7)}}{\partial G_j}\left((G_{39})^{**}\right) = s_{ij}$ | |
| Then taking into account equations and neglecting the terms of power 2, we obtain from | |
| $\frac{d\mathbb{G}_{36}}{dt} = -\left((a_{36}^{'})^{(7)} + (p_{36})^{(7)}\right)\mathbb{G}_{36} + (a_{36})^{(7)}\mathbb{G}_{37} - (q_{36})^{(7)}G_{36}^{*}\mathbb{T}_{37}$ | 573 |
| $\frac{d\mathbb{G}_{37}}{dt} = -\left((a_{37}^{'})^{(7)} + (p_{37})^{(7)}\right)\mathbb{G}_{37} + (a_{37})^{(7)}\mathbb{G}_{36} - (q_{37})^{(7)}G_{37}^{*}\mathbb{T}_{37}$ | 574 |
| $\frac{d\mathbb{G}_{38}}{dt} = -\left((a_{38}^{'})^{(7)} + (p_{38})^{(7)}\right)\mathbb{G}_{38} + (a_{38})^{(7)}\mathbb{G}_{37} - (q_{38})^{(7)}G_{38}^{*}\mathbb{T}_{37}$ | 575 |
| $\frac{d\mathbb{T}_{36}}{dt} = -\left((b_{36}^{'})^{(7)} - (r_{36})^{(7)}\right)\mathbb{T}_{36} + (b_{36})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(36)(j)}T_{36}^{*}\mathbb{G}_j\right)$ | 576 |
| $\frac{d\mathbb{T}_{37}}{dt} = -\left((b_{37}^{'})^{(7)} - (r_{37})^{(7)}\right)\mathbb{T}_{37} + (b_{37})^{(7)}\mathbb{T}_{36} + \sum_{j=36}^{38}\left(s_{(37)(j)}T_{37}^{*}\mathbb{G}_j\right)$ | 578 |
| $\frac{d\mathbb{T}_{38}}{dt} = -\left((b_{38}^{'})^{(7)} - (r_{38})^{(7)}\right)\mathbb{T}_{38} + (b_{38})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(38)(j)}T_{38}^{*}\mathbb{G}_j\right)$ | 579 |
| Obviously, these values represent an equilibrium solution<br><br>**ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 8:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(8)}$ *and* $(b_i^{''})^{(8)}$ Belong to $C^{(8)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^{*} + \mathbb{G}_i \quad , T_i = T_i^{*} + \mathbb{T}_i$<br><br>$\frac{\partial (a_{41}^{''})^{(8)}}{\partial T_{41}}(T_{41}^{*}) = (q_{41})^{(8)}$ , $\frac{\partial (b_i^{''})^{(8)}}{\partial G_j}\left((G_{43})^{*}\right) = s_{ij}$ | 580 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{40}}{dt} = -\left((a_{40}^{'})^{(8)} + (p_{40})^{(8)}\right)\mathbb{G}_{40} + (a_{40})^{(8)}\mathbb{G}_{41} - (q_{40})^{(8)}G_{40}^{*}\mathbb{T}_{41}$ | 581 |
| $\frac{d\mathbb{G}_{41}}{dt} = -\left((a_{41}^{'})^{(8)} + (p_{41})^{(8)}\right)\mathbb{G}_{41} + (a_{41})^{(8)}\mathbb{G}_{40} - (q_{41})^{(8)}G_{41}^{*}\mathbb{T}_{41}$ | 582 |
| $\frac{d\mathbb{G}_{42}}{dt} = -\left((a_{42}^{'})^{(8)} + (p_{42})^{(8)}\right)\mathbb{G}_{42} + (a_{42})^{(8)}\mathbb{G}_{41} - (q_{42})^{(8)}G_{42}^{*}\mathbb{T}_{41}$ | 583 |
| $\frac{d\mathbb{T}_{40}}{dt} = -\left((b_{40}^{'})^{(8)} - (r_{40})^{(8)}\right)\mathbb{T}_{40} + (b_{40})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(40)(j)}T_{40}^{*}\mathbb{G}_j\right)$ | 584 |

| | |
|---|---|
| $\frac{d\mathbb{T}_{41}}{dt} = -\left((b'_{41})^{(8)} - (r_{41})^{(8)}\right)\mathbb{T}_{41} + (b_{41})^{(8)}\mathbb{T}_{40} + \sum_{j=40}^{42}\left(s_{(41)(j)}T^*_{41}\mathbb{G}_j\right)$ | 585 |
| $\frac{d\mathbb{T}_{42}}{dt} = -\left((b'_{42})^{(8)} - (r_{42})^{(8)}\right)\mathbb{T}_{42} + (b_{42})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(42)(j)}T^*_{42}\mathbb{G}_j\right)$ | 586 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 9:** If the conditions of the previous theorem are satisfied and if the functions $(a_i'')^{(9)}$ *and* $(b_i'')^{(9)}$ Belong to $C^{(9)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br>**Proof: Denote** | 586 A |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial(a''_{45})^{(9)}}{\partial T_{45}}(T^*_{45}) = (q_{45})^{(9)} \ , \ \frac{\partial(b_i'')^{(9)}}{\partial G_j}((G_{47})^*) = s_{ij}$ | |
| Then taking into account equations 89 to 99 and neglecting the terms of power 2, we obtain from 99 to 44 | |
| $\frac{d\mathbb{G}_{44}}{dt} = -\left((a'_{44})^{(9)} + (p_{44})^{(9)}\right)\mathbb{G}_{44} + (a_{44})^{(9)}\mathbb{G}_{45} - (q_{44})^{(9)}G^*_{44}\mathbb{T}_{45}$ | 586 B |
| $\frac{d\mathbb{G}_{45}}{dt} = -\left((a'_{45})^{(9)} + (p_{45})^{(9)}\right)\mathbb{G}_{45} + (a_{45})^{(9)}\mathbb{G}_{44} - (q_{45})^{(9)}G^*_{45}\mathbb{T}_{45}$ | 586 C |
| $\frac{d\mathbb{G}_{46}}{dt} = -\left((a'_{46})^{(9)} + (p_{46})^{(9)}\right)\mathbb{G}_{46} + (a_{46})^{(9)}\mathbb{G}_{45} - (q_{46})^{(9)}G^*_{46}\mathbb{T}_{45}$ | 586 D |
| $\frac{d\mathbb{T}_{44}}{dt} = -\left((b'_{44})^{(9)} - (r_{44})^{(9)}\right)\mathbb{T}_{44} + (b_{44})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(44)(j)}T^*_{44}\mathbb{G}_j\right)$ | 586 E |
| $\frac{d\mathbb{T}_{45}}{dt} = -\left((b'_{45})^{(9)} - (r_{45})^{(9)}\right)\mathbb{T}_{45} + (b_{45})^{(9)}\mathbb{T}_{44} + \sum_{j=44}^{46}\left(s_{(45)(j)}T^*_{45}\mathbb{G}_j\right)$ | 586 F |
| $\frac{d\mathbb{T}_{46}}{dt} = -\left((b'_{46})^{(9)} - (r_{46})^{(9)}\right)\mathbb{T}_{46} + (b_{46})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(46)(j)}T^*_{46}\mathbb{G}_j\right)$ | 586 G |
| **The characteristic equation of this system is** | 587 |
| $\left((\lambda)^{(1)} + (b'_{15})^{(1)} - (r_{15})^{(1)}\right)\{\left((\lambda)^{(1)} + (a'_{15})^{(1)} + (p_{15})^{(1)}\right)$<br><br>$\left[\left((\lambda)^{(1)} + (a'_{13})^{(1)} + (p_{13})^{(1)}\right)(q_{14})^{(1)}G^*_{14} + (a_{14})^{(1)}(q_{13})^{(1)}G^*_{13}\right)\Big]$<br><br>$\left((\lambda)^{(1)} + (b'_{13})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(14)}T^*_{14} + (b_{14})^{(1)}s_{(13),(14)}T^*_{14}\Big)$<br><br>$+\left((\lambda)^{(1)} + (a'_{14})^{(1)} + (p_{14})^{(1)}\right)(q_{13})^{(1)}G^*_{13} + (a_{13})^{(1)}(q_{14})^{(1)}G^*_{14}\Big)$<br><br>$\left((\lambda)^{(1)} + (b'_{13})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(13)}T^*_{14} + (b_{14})^{(1)}s_{(13),(13)}T^*_{13}\Big)$<br><br>$\left((\lambda)^{(1)}\right)^2 + \left((a'_{13})^{(1)} + (a'_{14})^{(1)} + (p_{13})^{(1)} + (p_{14})^{(1)}\right)(\lambda)^{(1)}\Big)$ | |

$$\left((\lambda)^{(1)}\right)^2+\left((b_{13}')^{(1)}+(b_{14}')^{(1)}-(r_{13})^{(1)}+(r_{14})^{(1)}\right)(\lambda)^{(1)}\Big)$$

$$+\left(\left((\lambda)^{(1)}\right)^2+\left((a_{13}')^{(1)}+(a_{14}')^{(1)}+(p_{13})^{(1)}+(p_{14})^{(1)}\right)(\lambda)^{(1)}\right)(q_{15})^{(1)}G_{15}$$

$$+\left((\lambda)^{(1)}+(a_{13}')^{(1)}+(p_{13})^{(1)}\right)\left((a_{15})^{(1)}(q_{14})^{(1)}G_{14}^{*}+(a_{14})^{(1)}(a_{15})^{(1)}(q_{13})^{(1)}G_{13}^{*}\right)$$

$$\left((\lambda)^{(1)}+(b_{13}')^{(1)}-(r_{13})^{(1)}\right)s_{(14),(15)}T_{14}^{*}+(b_{14})^{(1)}s_{(13),(15)}T_{13}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(2)}+(b_{18}')^{(2)}-(r_{18})^{(2)}\}\{\left((\lambda)^{(2)}+(a_{18}')^{(2)}+(p_{18})^{(2)}\right)$$

$$\left[\left(\left((\lambda)^{(2)}+(a_{16}')^{(2)}+(p_{16})^{(2)}\right)(q_{17})^{(2)}\mathrm{G}_{17}^{*}+(a_{17})^{(2)}(q_{16})^{(2)}\mathrm{G}_{16}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(2)}+(b_{16}')^{(2)}-(r_{16})^{(2)}\right)s_{(17),(17)}\mathrm{T}_{17}^{*}+(b_{17})^{(2)}s_{(16),(17)}\mathrm{T}_{17}^{*}\right)$$

$$+\left(\left((\lambda)^{(2)}+(a_{17}')^{(2)}+(p_{17})^{(2)}\right)(q_{16})^{(2)}\mathrm{G}_{16}^{*}+(a_{16})^{(2)}(q_{17})^{(2)}\mathrm{G}_{17}^{*}\right)$$

$$\left(\left((\lambda)^{(2)}+(b_{16}')^{(2)}-(r_{16})^{(2)}\right)s_{(17),(16)}\mathrm{T}_{17}^{*}+(b_{17})^{(2)}s_{(16),(16)}\mathrm{T}_{16}^{*}\right)$$

$$\left(\left((\lambda)^{(2)}\right)^2+\left((a_{16}')^{(2)}+(a_{17}')^{(2)}+(p_{16})^{(2)}+(p_{17})^{(2)}\right)(\lambda)^{(2)}\right)$$

$$\left(\left((\lambda)^{(2)}\right)^2+\left((b_{16}')^{(2)}+(b_{17}')^{(2)}-(r_{16})^{(2)}+(r_{17})^{(2)}\right)(\lambda)^{(2)}\right)$$

$$+\left(\left((\lambda)^{(2)}\right)^2+\left((a_{16}')^{(2)}+(a_{17}')^{(2)}+(p_{16})^{(2)}+(p_{17})^{(2)}\right)(\lambda)^{(2)}\right)(q_{18})^{(2)}\mathrm{G}_{18}$$

$$+\left((\lambda)^{(2)}+(a_{16}')^{(2)}+(p_{16})^{(2)}\right)\left((a_{18})^{(2)}(q_{17})^{(2)}\mathrm{G}_{17}^{*}+(a_{17})^{(2)}(a_{18})^{(2)}(q_{16})^{(2)}\mathrm{G}_{16}^{*}\right)$$

$$\left((\lambda)^{(2)}+(b_{16}')^{(2)}-(r_{16})^{(2)}\right)s_{(17),(18)}\mathrm{T}_{17}^{*}+(b_{17})^{(2)}s_{(16),(18)}\mathrm{T}_{16}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(3)}+(b_{22}')^{(3)}-(r_{22})^{(3)}\}\{\left((\lambda)^{(3)}+(a_{22}')^{(3)}+(p_{22})^{(3)}\right)$$

$$\left[\left(\left((\lambda)^{(3)}+(a_{20}')^{(3)}+(p_{20})^{(3)}\right)(q_{21})^{(3)}G_{21}^{*}+(a_{21})^{(3)}(q_{20})^{(3)}G_{20}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(21)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(21)}T_{21}^{*}\right)$$

$$+\left(\left((\lambda)^{(3)}+(a_{21}')^{(3)}+(p_{21})^{(3)}\right)(q_{20})^{(3)}G_{20}^{*}+(a_{20})^{(3)}(q_{21})^{(1)}G_{21}^{*}\right)$$

$$\left(\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(20)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(20)}T_{20}^{*}\right)$$

$$\left(\left((\lambda)^{(3)}\right)^2+\left((a_{20}')^{(3)}+(a_{21}')^{(3)}+(p_{20})^{(3)}+(p_{21})^{(3)}\right)(\lambda)^{(3)}\right)$$

$$\left(\left((\lambda)^{(3)}\right)^2+\left((b_{20}')^{(3)}+(b_{21}')^{(3)}-(r_{20})^{(3)}+(r_{21})^{(3)}\right)(\lambda)^{(3)}\right)$$

$$+\left(\left((\lambda)^{(3)}\right)^2+\left((a_{20}')^{(3)}+(a_{21}')^{(3)}+(p_{20})^{(3)}+(p_{21})^{(3)}\right)(\lambda)^{(3)}\right)(q_{22})^{(3)}G_{22}$$

$$+\left((\lambda)^{(3)}+(a_{20}')^{(3)}+(p_{20})^{(3)}\right)\left((a_{22})^{(3)}(q_{21})^{(3)}G_{21}^{*}+(a_{21})^{(3)}(a_{22})^{(3)}(q_{20})^{(3)}G_{20}^{*}\right)$$

$$\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(22)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(22)}T_{20}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(4)}+(b_{26}')^{(4)}-(r_{26})^{(4)}\{\left((\lambda)^{(4)}+(a_{26}')^{(4)}+(p_{26})^{(4)}\right)$$

$$\left[\left(\left((\lambda)^{(4)}+(a_{24}')^{(4)}+(p_{24})^{(4)}\right)(q_{25})^{(4)}G_{25}^{*}+(a_{25})^{(4)}(q_{24})^{(4)}G_{24}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(4)}+(b_{24}')^{(4)}-(r_{24})^{(4)}\right)s_{(25),(25)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(25)}T_{25}^{*}\right)$$

$$+\left(\left((\lambda)^{(4)}+(a_{25}')^{(4)}+(p_{25})^{(4)}\right)(q_{24})^{(4)}G_{24}^{*}+(a_{24})^{(4)}(q_{25})^{(4)}G_{25}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}+(b_{24}')^{(4)}-(r_{24})^{(4)}\right)s_{(25),(24)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(24)}T_{24}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}\right)^2+\left((a_{24}')^{(4)}+(a_{25}')^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)$$

$$\left(\left((\lambda)^{(4)}\right)^2+\left((b_{24}')^{(4)}+(b_{25}')^{(4)}-(r_{24})^{(4)}+(r_{25})^{(4)}\right)(\lambda)^{(4)}\right)$$

$$+\left(\left((\lambda)^{(4)}\right)^2+\left((a_{24}')^{(4)}+(a_{25}')^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)(q_{26})^{(4)}G_{26}$$

$$+\left((\lambda)^{(4)}+(a_{24}')^{(4)}+(p_{24})^{(4)}\right)\left((a_{26})^{(4)}(q_{25})^{(4)}G_{25}^{*}+(a_{25})^{(4)}(a_{26})^{(4)}(q_{24})^{(4)}G_{24}^{*}\right)$$

$$\left((\lambda)^{(4)}+(b_{24}')^{(4)}-(r_{24})^{(4)}\right)s_{(25),(26)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(26)}T_{24}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(5)}+(b_{30}')^{(5)}-(r_{30})^{(5)}\{\left((\lambda)^{(5)}+(a_{30}')^{(5)}+(p_{30})^{(5)}\right)$$

$$\left[\left(\left((\lambda)^{(5)}+(a_{28}')^{(5)}+(p_{28})^{(5)}\right)(q_{29})^{(5)}G_{29}^{*}+(a_{29})^{(5)}(q_{28})^{(5)}G_{28}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(5)}+(b_{28}')^{(5)}-(r_{28})^{(5)}\right)s_{(29),(29)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(29)}T_{29}^{*}\right)$$

$$+\left(\left((\lambda)^{(5)}+(a_{29}')^{(5)}+(p_{29})^{(5)}\right)(q_{28})^{(5)}G_{28}^{*}+(a_{28})^{(5)}(q_{29})^{(5)}G_{29}^{*}\right)$$

$$\left(\left((\lambda)^{(5)}+(b_{28}')^{(5)}-(r_{28})^{(5)}\right)s_{(29),(28)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(28)}T_{28}^{*}\right)$$

$$\left(\left((\lambda)^{(5)}\right)^2+\left((a_{28}')^{(5)}+(a_{29}')^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)$$

$$\left((\lambda)^{(5)}\right)^2+\left((b_{28}')^{(5)}+(b_{29}')^{(5)}-(r_{28})^{(5)}+(r_{29})^{(5)}\right)(\lambda)^{(5)}\Big)$$

$$+\left(\left((\lambda)^{(5)}\right)^2+\left((a_{28}')^{(5)}+(a_{29}')^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)(q_{30})^{(5)}G_{30}$$

$$+\left((\lambda)^{(5)}+(a_{28}')^{(5)}+(p_{28})^{(5)}\right)\left((a_{30})^{(5)}(q_{29})^{(5)}G_{29}^{*}+(a_{29})^{(5)}(a_{30})^{(5)}(q_{28})^{(5)}G_{28}^{*}\right)$$

$$\left((\lambda)^{(5)}+(b_{28}')^{(5)}-(r_{28})^{(5)}\right)s_{(29),(30)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(30)}T_{28}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(6)}+(b_{34}')^{(6)}-(r_{34})^{(6)}\{\left((\lambda)^{(6)}+(a_{34}')^{(6)}+(p_{34})^{(6)}\right)$$

$$\left[\left((\lambda)^{(6)}+(a_{32}')^{(6)}+(p_{32})^{(6)}\right)(q_{33})^{(6)}G_{33}^{*}+(a_{33})^{(6)}(q_{32})^{(6)}G_{32}^{*}\right)\Big]$$

$$\left((\lambda)^{(6)}+(b_{32}')^{(6)}-(r_{32})^{(6)}\right)s_{(33),(33)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(33)}T_{33}^{*}\Big)$$

$$+\left((\lambda)^{(6)}+(a_{33}')^{(6)}+(p_{33})^{(6)}\right)(q_{32})^{(6)}G_{32}^{*}+(a_{32})^{(6)}(q_{33})^{(6)}G_{33}^{*}\Big)$$

$$\left((\lambda)^{(6)}+(b_{32}')^{(6)}-(r_{32})^{(6)}\right)s_{(33),(32)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(32)}T_{32}^{*}\Big)$$

$$\left((\lambda)^{(6)}\right)^2+\left((a_{32}')^{(6)}+(a_{33}')^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\Big)$$

$$\left((\lambda)^{(6)}\right)^2+\left((b_{32}')^{(6)}+(b_{33}')^{(6)}-(r_{32})^{(6)}+(r_{33})^{(6)}\right)(\lambda)^{(6)}\Big)$$

$$+\left(\left((\lambda)^{(6)}\right)^2+\left((a_{32}')^{(6)}+(a_{33}')^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\right)(q_{34})^{(6)}G_{34}$$

$$+\left((\lambda)^{(6)}+(a_{32}')^{(6)}+(p_{32})^{(6)}\right)\left((a_{34})^{(6)}(q_{33})^{(6)}G_{33}^{*}+(a_{33})^{(6)}(a_{34})^{(6)}(q_{32})^{(6)}G_{32}^{*}\right)$$

$$\left((\lambda)^{(6)}+(b_{32}')^{(6)}-(r_{32})^{(6)}\right)s_{(33),(34)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(34)}T_{32}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(7)}+(b_{38}')^{(7)}-(r_{38})^{(7)}\{\left((\lambda)^{(7)}+(a_{38}')^{(7)}+(p_{38})^{(7)}\right)$$

$$\left[\left((\lambda)^{(7)}+(a_{36}')^{(7)}+(p_{36})^{(7)}\right)(q_{37})^{(7)}G_{37}^{*}+(a_{37})^{(7)}(q_{36})^{(7)}G_{36}^{*}\right)\Big]$$

$$\left((\lambda)^{(7)}+(b_{36}')^{(7)}-(r_{36})^{(7)}\right)s_{(37),(37)}T_{37}^{*}+(b_{37})^{(7)}s_{(36),(37)}T_{37}^{*}\Big)$$

$$+\left((\lambda)^{(7)}+(a_{37}')^{(7)}+(p_{37})^{(7)}\right)(q_{36})^{(7)}G_{36}^{*}+(a_{36})^{(7)}(q_{37})^{(7)}G_{37}^{*}\Big)$$

$$\left((\lambda)^{(7)}+(b_{36}')^{(7)}-(r_{36})^{(7)}\right)s_{(37),(36)}T_{37}^{*}+(b_{37})^{(7)}s_{(36),(36)}T_{36}^{*}\Big)$$

$$\left((\lambda)^{(7)}\right)^2+\left((a_{36}')^{(7)}+(a_{37}')^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\Big)$$

$$\left((\lambda)^{(7)}\right)^2+\left((b_{36}')^{(7)}+(b_{37}')^{(7)}-(r_{36})^{(7)}+(r_{37})^{(7)}\right)(\lambda)^{(7)}\Big)$$

$$+\left(\left((\lambda)^{(7)}\right)^2+\left((a_{36}')^{(7)}+(a_{37}')^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\right)(q_{38})^{(7)}G_{38}$$

$$+\left((\lambda)^{(7)}+(a_{36}')^{(7)}+(p_{36})^{(7)}\right)\left((a_{38})^{(7)}(q_{37})^{(7)}G_{37}^{*}+(a_{37})^{(7)}(a_{38})^{(7)}(q_{36})^{(7)}G_{36}^{*}\right)$$

$$\left((\lambda)^{(7)}+(b_{36}')^{(7)}-(r_{36})^{(7)}\right)s_{(37),(38)}T_{37}^{*}+(b_{37})^{(7)}s_{(36),(38)}T_{36}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(8)}+(b_{42}')^{(8)}-(r_{42})^{(8)}\}\{\left((\lambda)^{(8)}+(a_{42}')^{(8)}+(p_{42})^{(8)}\right)$$

$$\left[\left(\left((\lambda)^{(8)}+(a_{40}')^{(8)}+(p_{40})^{(8)}\right)(q_{41})^{(8)}G_{41}^{*}+(a_{41})^{(8)}(q_{40})^{(8)}G_{40}^{*}\right)\right]$$

$$\left((\lambda)^{(8)}+(b_{40}')^{(8)}-(r_{40})^{(8)}\right)s_{(41),(41)}T_{41}^{*}+(b_{41})^{(8)}s_{(40),(41)}T_{41}^{*}\Big)$$

$$+\left(\left((\lambda)^{(8)}+(a_{41}')^{(8)}+(p_{41})^{(8)}\right)(q_{40})^{(8)}G_{40}^{*}+(a_{40})^{(8)}(q_{41})^{(8)}G_{41}^{*}\right)$$

$$\left((\lambda)^{(8)}+(b_{40}')^{(8)}-(r_{40})^{(8)}\right)s_{(41),(40)}T_{41}^{*}+(b_{41})^{(8)}s_{(40),(40)}T_{40}^{*}\Big)$$

$$\left((\lambda)^{(8)}\right)^2+\left((a_{40}')^{(8)}+(a_{41}')^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\Big)$$

$$\left((\lambda)^{(8)}\right)^2+\left((b_{40}')^{(8)}+(b_{41}')^{(8)}-(r_{40})^{(8)}+(r_{41})^{(8)}\right)(\lambda)^{(8)}\Big)$$

$$+\left(\left((\lambda)^{(8)}\right)^2+\left((a_{40}')^{(8)}+(a_{41}')^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\right)(q_{42})^{(8)}G_{42}$$

$$+\left((\lambda)^{(8)}+(a_{40}')^{(8)}+(p_{40})^{(8)}\right)\left((a_{42})^{(8)}(q_{41})^{(8)}G_{41}^{*}+(a_{41})^{(8)}(a_{42})^{(8)}(q_{40})^{(8)}G_{40}^{*}\right)$$

$$\left((\lambda)^{(8)}+(b_{40}')^{(8)}-(r_{40})^{(8)}\right)s_{(41),(42)}T_{41}^{*}+(b_{41})^{(8)}s_{(40),(42)}T_{40}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(9)}+(b_{46}')^{(9)}-(r_{46})^{(9)}\}\{\left((\lambda)^{(9)}+(a_{46}')^{(9)}+(p_{46})^{(9)}\right)$$

$$\left[\left(\left((\lambda)^{(9)}+(a_{44}')^{(9)}+(p_{44})^{(9)}\right)(q_{45})^{(9)}G_{45}^{*}+(a_{45})^{(9)}(q_{44})^{(9)}G_{44}^{*}\right)\right]$$

$$\left((\lambda)^{(9)}+(b_{44}')^{(9)}-(r_{44})^{(9)}\right)s_{(45),(45)}T_{45}^{*}+(b_{45})^{(9)}s_{(44),(45)}T_{45}^{*}\Big)$$

$$+\left(\left((\lambda)^{(9)}+(a_{45}')^{(9)}+(p_{45})^{(9)}\right)(q_{44})^{(9)}G_{44}^{*}+(a_{44})^{(9)}(q_{45})^{(9)}G_{45}^{*}\right)$$

$$\left((\lambda)^{(9)}+(b_{44}')^{(9)}-(r_{44})^{(9)}\right)s_{(45),(44)}T_{45}^{*}+(b_{45})^{(9)}s_{(44),(44)}T_{44}^{*}\Big)$$

$$\left((\lambda)^{(9)}\right)^2+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\Big)$$

$$\left((\lambda)^{(9)}\right)^2+\left((b'_{44})^{(9)}+(b'_{45})^{(9)}-(r_{44})^{(9)}+(r_{45})^{(9)}\right)(\lambda)^{(9)}\Big)$$

$$+\left(\left((\lambda)^{(9)}\right)^2+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\right)(q_{46})^{(9)}G_{46}$$

$$+\left((\lambda)^{(9)}+(a'_{44})^{(9)}+(p_{44})^{(9)}\right)\left((a_{46})^{(9)}(q_{45})^{(9)}G^*_{45}+(a_{45})^{(9)}(a_{46})^{(9)}(q_{44})^{(9)}G^*_{44}\right)$$

$$\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(46)}T^*_{45}+(b_{45})^{(9)}s_{(44),(46)}T^*_{44}\Big)\}=0$$

**And as one sees, all the coefficients are positive. It follows that all the roots have negative real part, and this proves the theorem.**

Note: **Same dovetailing explanation holds good for d/dt (partial differential with respect to t), d/dt,d^2/dt^2 (acceralation: double dot). Such an exposition is helpful in optimisation problems, duality of motion, wave motion a homogeneous functions studies, and control theory which we intend to incorporate in future.** It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## Persistence, Coexistence And Thrusty Permanence For Interacting Structured Particles, Transactions In Finance, Econometric Productivity Functions, Environment, Interfaces, Cosmological Interactions Et Al

## Josef Hofbauer And Sebastian J. Schreiber Model Thereof: A Functional Appraisal

ABSTRACT. **JOSEF HOFBAUER AND SEBASTIAN J. SCHREIBER** studied the dynamics of interacting structured systems (in their case population dynamics) can be modeled by $\frac{dxi}{dt}= A_i (x) xi$ where $x_i \epsilon R^n_{i,} x = (x_1,\ldots,x_k)$, and $A_i(x)$ are matrices with non-negative off- diagonal entries. These models are permanent if there exists

a positive global attractor and are robustly permanent if they remain permanent following perturbations of $A_i(x)$. Necessary and sufficient conditions for robust permanence are derived using dominant Lyapunov exponents $X_i(\mu)$ of the $A_i(x)$ with respect to invariant measures $\mu$. The necessary conditions requires $\max_i X_i(\mu)>0$ for all ergodic measures with support in the boundary of the non-negative cone. The sufficient condition requires that the boundary admits a Morse decomposition such that $\max_i Xi(\mu)>0$ for all invariant measures $\mu$ supported by a component of the Morse decomposition. When the Morse components are Axiom A, uniquely ergodic, or support all but one population, the necessary and sufficient conditions are equivalent. Applications to spatial ecology, epidemiology, and gene networks are given. Appeared in *Journal of Differential Equations*, 248, 1955-1971 (2010)
It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## INTRODUCTION

A fundamental issue in population biology is what are the minimal conditions to ensure the long-term survivorship for interacting populations whether they be viral particles, bio-chemicals, plants, or animals. When these conditions are met the interacting populations are said to persist or coexist. Since the pioneering work of Lotka and Volterra on competitive and predator prey interactions, Thompson, Nicholson, and Bailey on host-parasite interactions, and Kermack and McKendrick on disease outbreaks, nonlinear difference and differential equations have been used to understand conditions for population persistence [1-5]. One particularly important form of persistence for deterministic models is *permanence* or *uniform persistence* which corresponds to the existence of a global attractor bounded away from extinction of one or more species. When such an attractor exists, interacting populations are able to recover from "vigorous shake ups" of the of the population state. Permanence has been characterized from a topological perspective [6-8] and with average Lyapunov functions [9-12] for differential equation models of the form

$$\frac{dxi}{dt} = x_i f_i(x) \qquad i = 1,\ldots,k \tag{5}$$

Where $x_i$ is the abundance of population $i$, $x = (x_1,\ldots,x_k)$, and $f_i(x)$ is the per-capita growth of population $i$.

An important extension of the concept of the permanence is *robust permanence* [13-16] which requires that (1) remain permanent after sufficiently small perturbations of the per-capita growth rates $f_i$. The importance of robust permanence as one crude form of structural stability for population models are approximations to reality. Consequently, if nearby models (e.g. more realistic models) are not permanent despite the focal model being permanent, and then one can draw few conclusions about the persistence of the biological system being approximated by the model. One can view robust permanence as one crude form of structural stability for debits-credits particles, propositions, transactions in the celestial subterranean realm and ceratoid dualism, finance like debiting say a current account and crediting to another account, or recrediting to the account which had been earlier debited, further conduction of the cash credit account, black hole formation, guzzling of inventories for the output, creation of income from output, reinvestment of income to investment,gormadizing of matter by black holes, and emission of radiation therefore, gamma ray bursts, in fact most of the transactions and interactions in the universe are conservative and quantum mechanical universe, something like a account acting as predator and prey. More often than not as we have seen from Haimovici model such models could be applied to the problems, propositions, which are not subsumed by the concomitant problems in the bearing. the population models. For this perspective, it is not unexpected that there are permanent systems that cannot be approximated by robustly permanent systems [17]. In [14, 16], criteria for robust permanence were developed with respect to the average per-capita growth rates $\int f_i(x)d\mu(x)$ with respect to invariant probability measures $\mu$ supported on the boundary of the positive orthant (i.e. where one or more Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates are extinct). Roughly, these criteria for robust permanence require that the average per-capita growth rate (i.e. invasion rate) is positive for some missing species.

While (1) can account for many types of Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates. Interactions, it assumes that all individual within populations are exactly the same. However, theoretical biologists have long recognized that different individuals within a population may be in different states (e.g. different sizes or ages, living in different spatial locations) and these differences can have important consequences for population dynamics [18-24]. To account for how these differences influence persistence of interacting populations, we develop criteria for robust permanence for the dynamics of $k$ interacting structured populations. Roughly, these dynamics correspond to

replacing $x_i$ in (1) by vectors and $f_i(x)$ by matrices. In section 2, we describe these models in greater detail and review some basic concepts from dynamical systems theory. In section 3, we introduce the structured analog of average per-capita growth rates and define the concept of robustly unsaturated invariant sets. Necessary and sufficient conditions for an invariant set to be robustly unsaturated are proven. For Axiom A or uniquely ergodic invariant sets, the necessary and sufficient criteria are shown to be equivalent. In section 4, we use Morse decompositions and the criteria for robustly unsaturated invariant sets to develop necessary and sufficient conditions for robust permanence. In sections 5 and 6, we provide several applications of our results to spatially structured ecological models, structured epidemiological models, and models of gene networks. We apply the model to econometrics. To give an example investment needs inventory to be used towards the consummation of output, part of output is to be sold for consolidation of income, and some income might be replenished in the investment and so on. In finance a cash credit account could be debited or credited when there is say purchase of demand bills or discount of usance bills. Cash credit account is debited whenever cheques are issued towards the purchase of raw materials and so on. In fact some of the models have been applied to Organic psychiatry and we have shown how such models are useful in dealing with the propositions of Husserl, Derrida, Deleuze and others.

## 5. MODELS AND ASSUMPTIONS

Let $x_i$ denote the state of the $i$-th population that lies in the non-negative cone $C_i$ of $R^{ni.}$ Define $x=(x_1,\ldots x_k)\epsilon C$ to be the non-negative cone of $R^n$ where $n=\sum_{i=1}^{k} n_i$. Let $C_i^+$ denote positive cone { $x_i \epsilon\ C_i$: $\prod_i$i noitalupop rof$^j$

$i$ *and* $C^+=\prod_i C^+_i$ , the positive cone for the interacting populations. If $x \mapsto A_i\ (x)$ is a map into $n_i x n_i$ matrices that describes the growth of population I, then the dynamics of the interacting populations are given by

$$(2)\frac{dx_i}{dt} = A_i(x)x_i \ i=1,\ldots,k$$

Let $x.t$ denote the solution of (2), with the initial condition x, and more generally the semi flow generated by (2). Before stating our assumption on (2), recall a few definitions from dynamical systems. Given sets $I \subseteq R$ and $K \subseteq C$, let $K.I = \{x.t : t \epsilon I, x \epsilon K\}$. A set $K \subseteq C$ *is invariant* if K.t = K for all t>0, and *forward invariant* if $K.t \subset K$ for all $t > 0$. The *omega limit* set of a set $K \subseteq C$ equals $\omega(K) = \bigcap_{t\geq 0,} \overline{K.[t,\infty)}$. The *alpha limit* set of $K \subseteq C$ is $a(K) = \bigcap_{t\leq 0}$

$\overline{K.(-\infty,t]}$. Given a forward invariant set K, $B \subset K$ is an attractor for the semi flow *x.t restricted* to K provided there exists an open neighborhood $U \subseteq K$ of B such that $\omega(U) = B$. The stable set of a *compact invariant* set *K* is defined by

$$W^s(K) = \{x \in C : \omega(x) \neq \emptyset \text{ and } \omega(x) \subseteq K\}.$$

The semi flow generated by (2) is *dissipative* if there exists an attractor B with $W^s(B) = C$.

Throughout this paper, we make the following assumptions:

**A1**: $x \mapsto A_i(x)$ is continuous

**A2:** $A_i(x)$ is irreducible and has non-negative off- diagonal entries,

**A3:** *x.t* is defined for all $t \geq 0$

**A4:** *x.t* is dissipative with global attractor $\Gamma(A)$.

A1 is a basic regularity assumption. The non-negative of off- diagonal entries in A2 implies that there are no negative feedbacks between individuals of different states in population *i*. This assumption is meet for many types of structured models as discussed in sections 5 and 6. The irreducibility assumption A2 is generally meet and implies that all individuals within a population can pass through all states. In Remark1, we discuss how this irreducibility assumption can be relaxed. A3 ensures that the population dynamics are defined for all future time. A4 requires that population densities/ abundances eventually are uniformly bounded, a condition that should be meet for an biologically realistic model.

We say that the semi flow of (2) is *permanent* if this semi flow is dissipative and there is a positive attractor $B \subset C^+$ such that $W^s(B) \subset C^+$. We will denote this positive attractor by $\Lambda(A)$. Permanence ensures that populations can recover from rare large perturbations and allows for a diversity of dynamical behaviors.

For K $\subset$ C and $\delta > 0$, define the $\delta$- *neighborhood* of K as

$$N\delta(K) = \{x \in C : \backslash x - y < \delta \text{ for some } y \in K\}.$$

We define a $\delta$- *perturbation of (2)* to be a system of the form

$$\frac{dx}{dt} = \tilde{A}(x)x$$

that satisfies assumptions A1-A4, $\Gamma(\tilde{A}) \subset N$ *( $\Gamma(A)$), and* $||A(x) - \tilde{A}(x)|| \leq \delta$ *for all* $x \in N$ *($\Gamma(A)$).* (2) is *robustly permanent* if there exist >0 and $\epsilon > 0$ such that all $\delta$- perturbations of (2) are permanent, and $d(\mathcal{A}(\tilde{A}), C \backslash C^{+}) > \epsilon$ for all $\delta$- perturbations i.e., there is a common/ uniform region of repulsion around the boundary.

## INUNDATION OF COMPACT SETS

We begin by studying the linear skew product flows on $\Gamma(A) \times R^{n}_{i}$ defined by $(x,y).t = (x.t, B_i(t,x)y)$ where $Y(t) = B_i(t,x)$ is the solution to $\frac{dY}{dt}(t) = A_i(x.t)Y(t)$ with y(0) equal to the identify matrix. Our assumption that $A_i$ is irreducible with non- negative off diagonal entries implies that $B_i$ (t,x)$C_i \subset C_i^{+}$ for all $x$ and t > 0 (see, e.g., [25] ). A result of Ruelle [26, Prop, 3.2] implies that there exist continuous maps $u_i\, v_i : \Gamma(A) \rightarrow C_i^{+}$ with $u_i(x)| = |v_i(x)| = 1$ such that The line bundle $E_i(x)$ spanned by $u_i(x)$ is invariant i.e. $E_i(x.t) = B_i(x,t)E_i(x)$ for all t ≥ 0.The vector bundle $F_i(x)$ perpendicular to $v_i(x)$ is invariant i.e. $F_i(x.t) = B_i(x,t)F_i(x)$ for all $t \geq 0$.There exist constants $a > 0$ *and* $\beta > 0$ such that(3) $||B_i(t,x)F_i(x)|| \leq a\, exp(-\beta t)||B_i(t,x)E_i(x)||$ For all $x \in K$ *and* $t \geq 0$.**Interaction rates.** Given any x $\epsilon$ $\Gamma$(A), define the *interaction and inundation rate* of *species i at particle/proposition/problem/ state x as* $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right) in\ ||B_i(t,x)||$ when $x_i$ =0, $\lambda_i(x)$ provides an upper bound to the rate of growth of population *i* when introduced at infinitesimally small densities. Important properties of this invasion rate are summarized in the following two propositions. For instance, the second proposition implies if $(x.t)_i$ stays bounded away from zero (i.e. population *i* persist), then

$\lambda_i(x) = 0$. Alternatively, if $\lambda_i(x) < 0$,then population *i* is doomed to extinction. **Proposition 1.** $\lambda_i(x)$ *satisfies the following properties:* $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right) In\ |B_i\ (t,x)v|$ *for any* $v > 0$ *in* $R^{ni}$, *and* $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right) \int^{t}_{0}(A_i(x.s)u_i(x.s), u_i(x.s))ds.$

**Proof:** To prove the first property, we first show that

$$\lambda i(x) = \underset{t\to\infty}{limsup}\frac{1}{t} In\ |B_i(t,x)u_i(x)|.$$

To this end, let $v \in R^{ni}$ be any non-zero vector. Since $R^{ni} = E_{i(x)}$ $F_i(x)$, we can write $v = au_i(x) + w$ with $a \in R$ and $w \in F_i(x)$. Equation (3) implies

$|B_i(t,x)v| \leq a|B_i(t,x)u_i(x)| + |B_i(t,x)w| \leq |B_i(t,x)u_i(x)|(a + \alpha + exp(-\beta t)|w|)$. Hence, $\limsup_{t\to\infty} \frac{1}{t} In\, B_i(t,x)\, v| \leq$

$$\limsup_{t\to\infty} \frac{1}{t} In\, B_i(t,x)\, u_i(x) \leq \lambda_i(x)$$

for all non zero vectors $v \in R^{ni}$. Since $||B_i(t,x)|| = sup|v|=1 B_i(t,x)v|$, this inequality implies that $\lambda i(x) = \limsup_{t\to\infty} \frac{1}{t}$ In $|B_i(t,x)v|$, this inequality implies that $\lambda_i(x) = limsup_{t\to\infty} \left(\frac{1}{t}\right)$ In $|t,x)u_i(x)|$. Now let v $\epsilon$ $C_i^+$. Then, we can write $v = au_i(x) + w$ with $a > 0$ and $w \in F_i$.

Equation (3) implies

$$|B_i(t,x)v| \geq a|B_i(t,x)u_i(x)| - |B_i(t,x)w|$$

$$\geq |B_i(t,x)u_i(x)|(a - \alpha + exp(-\beta t)|w|).$$

Since a>0, this inequality implies that

$$\lambda i(x) \geq \limsup_{t\to\infty} \frac{1}{t} \text{In } |B_i(t,x)v| \geq \lambda i(x) \geq \limsup_{t\to\infty} \frac{1}{t} \text{In } |B_i(t,x)u_i(x)| = \lambda i(x)$$

To prove the second assertion, let $b_i(t,x) = In\, |B_i(t,x)u_i(x)|$. Invariance of $u_i(x)$ implies that $b_i(t,x)$ is additive:

$$B_i(t+s,x) = In|B_i(t+s,x)\, u_i(x)| = In\, |B_i(s,x.t)B_i(t,x)u_i(x)$$

$$= \text{In } |B_i(s,x.t)\, u_i(x.t)||\, B_i(t,x)\, u_i(x)|$$

$$= \text{In } b_i(s,t.x)\, b_i(t,x)$$

Additivity of $b_i(t,x)$ and the fact that $b_i(0,x) = 0$ implies

$$\frac{d}{dt} b_i(t,x) = \lim_{s\to 0} \frac{bi(t+s,x - bi(t,x)}{S-}$$

$$= \lim_{s\to 0} \frac{bi(s,x.t)}{S} = \frac{d}{ds}|_{s=0}\, b_i(s,x.t)$$

$$=$$

$$= (A_i(x.t)u_i(x.s), u_i(x.s),$$

The Fundamental Theorem of Calculus implies

$$B_i(t,x)= \int_0^t (A_i(x.s)\, u_i(x.s), u_i(x.s))ds\backslash$$

and the second assertion follows.

**Proposition 2:** *For the solutions of (2) we have:* if $x_i>0$, then $\lambda i(x) \leq 0$. If $\lambda i(x)<0$, then $\lim_{t\to\infty} (x.t)_i=0$ If $limsup_{t\to\infty}$ $(x.t)_i|>0$, then $\lambda i(x)=0$.

**Proof:** First, assume that x satisfies $x_i>0$. Since the semi flow is dissipative, there exists $\gamma >0$ such that $|x.t| \leq \gamma$ for all $t \geq 0$. Proposition 1 and the definition of the skew product flow imply

$$\lambda i(x) = \limsup_{t\to\infty} \frac{1}{t} \text{In } |B_i(t,x)x_i| = \limsup_{t\to\infty} \frac{1}{t} \text{In } | (x.t)_i |$$

$$\leq \limsup_{t\to\infty} \frac{In\,\gamma}{t} = 0$$

To prove the second assertion, assume that $\lambda i(x) < 0$. $If\ x_i=0$, then the invariance of the faces of C imply that $(x.t)_i=0$ for all $t\geq 0$. Alternatively, if $x_i> 0$, then Proposition I and the skew product flow imply

$$\limsup_{t\to\infty} \frac{1}{t} In|(x.t)_i| = \lambda i(x) < 0.$$

The final assertion follows from the first two assertions. 5

**3.2. Invariant measures.** We review some definitions from ergodic theory. Given a Borel probability measure $\mu$ on ***C***, the *support* of $\mu$, denoted supp( $\mu$),is the smallest closed set whose complement has measure zero. A Borel probability measure $\mu$ is called *invariant* for (2) provided that Integral $h(x)d\mu(x) =$ Integral $h(x.t)d\,\mu(x)$for all $t \geq 0$ and for all bounded continuous functions $h$:C→R. An invariant measure $\mu$ is called *ergodic* provided that $\mu(B)$=0 or 1 for any invariant Borel set *B.*

For an invariant measure $\mu$ define the invasion rate species $i$ with respect to $\mu$ as

$$\lambda_i(\mu) = \int_c \langle A_i(x)u_i(x), u_i(x)\rangle d\mu(x)$$

**Proposition 3:** Let $\mu$ be an *invariant measure.* $\lambda_i(\mu)$ *Satisfies the following properties:* $\lambda_i(x) = \lim_{t\to\infty} \frac{1}{t}\int_0^t \langle A_i(x.s)u_i(x.s), u_i(x.s)\rangle, ds$ exists for $\mu$ - *almost every* $x$. *Moreover, if* $\mu$ *is ergodic,* $\lambda_i(x) = \lambda_i(\mu)$ for $\mu$-*almost every* $x$.*If* $\mu$ *is ergodic, then there exists* $I \subset \{1, \dots, k\}$such that $\mu(\Pi_{i\in I} C_i^+)$=1 and $\lambda_i(\mu) = 0$ *for all* $i \in I$.

**Proof:** Let $\mu$ be an invariant measure. Define $h_i(x) = \langle A_i(x)u_i(x), u_i(x)\rangle$. The Birkhoff's' ergodic theorem and Proposition 1 imply

$$\lambda_i(x) = \lim_{t\to\infty} \frac{1}{t}\int_0^t h_i(x.s)ds$$

Exists for $\mu$- almost every $x$. Moreover, $\lambda_i(x) = \lambda_i(\mu)$ $\mu$- almost surely if $\mu$ is ergodic.

Assume $\mu$ is ergodic. By ergodicity and invariance of the faces of **C**, there exists $I \subset \{1, \dots, k\}$ such that $\mu(\Pi_{i\in I} C_i^+)$=1. Let $K \subset \Pi_{i\in I} C_i^+$ be a compact set such that $\mu(K) > 0$. The Poincare recurrence theorem and the Birkhoff ergodic theorem imply that there is $x \in \Pi_{i\in I} C_i^+$such that $x.t_n \in K$ for some $t_n \uparrow \infty$ and $\lambda_i(x) = \lambda_i(\mu)$ for all $i$. The third assertion of Proposition 2 implies $\lambda_i(\mu) = 0$ for all $i \in I$.

**3.3. Robustly unsaturated sets**. Let $K \subset \partial C$ be a compact isolated invariant set for the flow $x.t$ restricted to $\partial C$.*K* is unsaturated if $W^s(K) \subset \partial C$ and K is isolated for $x.t$. If *K* is not unsaturated, then *K* is saturated. *K* is robustly unsaturated for (2) if there exists $\delta > 0$ such that the continuation of K for any $\delta -$ perturbation of (2) is unsaturated.**Theorem 1:** Let *K* be a compact isolated invariant set for $x.t$ restricted to$\partial C$. If one of the following equivalent conditions hold For all invariant measures $\mu$ supported by *K* (4) $\max_{1\le i\le k} \lambda_i(\mu) = 0$

There exist $p_i > 0$ such that (5) $\sum_{1\le i\le k} p_i \lambda_i(\mu) = 0$

For all ergodic probability measures supported by *K*.

then *K* is robustly unsaturated for (2). Alternatively, if $x \to A(x)$ is twice continuously differentiable and *K* is robustly unsaturated, then (4) holds for all ergodic measures $\mu$ supported by *K*.

**Remark 1**: For some applications (e.g., the disease model considered in section 6.1) it useful to relax the irreducibility assumption A2. For instance, if there exists an open neighborhood U of K such that for each $i$, $A_i(x)$ has a fixed off-diagonal sign pattern for all $x \in U$,then $A_i(x)$ can be decomposed into a finite number, say $m_i$, of

irreducible components. For each of these irreducible components, one can define $\lambda_i^j(\mu) = \int_C \langle A_i^j(x)u_i^j(x), u_i^j(x)\rangle \, d\mu(x)$ where $A_i^j(x)$ is the submatrix of $A_i(x)$ corresponding to the $j$-th irreducible component of $A_i(x)$ and $u_i^j(x)$ is the continuous invariant sub bundle for the irreducible linear cocycle determined by $A_i^j(x)$. If we define $\lambda_i(\mu) = \max_{1\le j\le m_i} \lambda_i^j(\mu)$, then all of the assertions of Theorem 1 still hold.

**Proof:** To prove the first assertion of the theorem, we need the following lemma. We call an invariant measure $\mu$ for (2) saturated if $\lambda_i(\mu) \le 0$ for all $i$

**Lemma 1**: Let $\delta_n$ be a non –negative sequence that converges to zero as $n\to\infty$. If $\mu_n$ are saturated invariant measures for $\delta_n -$ perturbations of (2), then the weak$^*$ limit points of $\{\mu_n\}_{n=1}^{\infty}$ is a non-empty set consisting of saturated invariant measures for (2).

**Proof:** Let $\delta_n$ be a non-negative sequence that coverages to zero. Let $\mu_n$ be saturated invariant measures for $\delta_n$ perturbations, $\dot{x} = A^n(x)x$, of $\dot{x} = A(x)x$. By weak$^*$ compactness of Borel probability measures supported on $N_1(\Gamma(A))$, there exist weak$^*$ limit points of $\mu_n$. Let $\mu$ be such a weak$^*$ limit point. Since $\mu_n$ are supported by $N_1(\Gamma(A))$, for all $n$ sufficiently large, $\mu$ is supported by $N_1(\Gamma(A))$.To verify that $\mu$ is an invariant measure for (2),let $h: C \to R$ be a bounded continuous function. Let $t > 0$ and $\in > 0$ be given. Choose $n$ sufficiently large so that $|\int_c h(x)\, d\mu(x) - \int_c h(x) d\mu_n(x)| \le \in$, $|\int_c h(x.t)\, d\mu(x) - \int_c h(x.t) d\mu_n(x)| \le \in$, and $|h(x.t) - h(x^n.t)| \le \in$ for all $x \in N_1(\Gamma(A))$.Then $|\int_c h(x.t) - h(x) d\mu(x)|$ is

$\le |\int_c h(x.t)\, d\mu(x) - \int_c h(x.t)\, d\mu_n(x)| + |\int_c h(x.t)\, d\mu_n(x) - \int_c h(x)\, d\mu(x)|$

$\le \in + \int_c |h(x.t) - h(x^n.t)|\, d\mu_n(x)| + |\int_c h(x^n.t)|\, d\mu_n(x) - \int_c h(x) d\mu(x)|$

$\le 2 \in + |\int_c h(x^n.t) - h(x)\, d\mu_n(x)| + |\int_c h(x) d\mu_n(x) - \int_c h(x) d\mu(x)|$

$\le 3 \in$

Where the last line follows from the invariance of $\mu_n$ for $x^n.t$. Since $\in > 0$ and t > 0 are arbitrary, we have $\int_c h(x.t)d\mu(x) = \int_c h(x)d\mu(x)$ for all t >0 and all bounded continuous functions $h: C \to R$. It follows that $\mu$ is an invariant measure for (2). To see that $\mu$ is saturated, define $h_i^n(x) = \langle A_i^n(x)u_i^n(x), u_i^n(x)\rangle$ where $u_i^n(x)$ spans the invariant one-dimensional bundle given by $[26, Prop.3.2]$ for $\frac{dx}{dt} = A^n(x)x$. Since $A_i^n(x) \to A_i(x)$ uniformly for $x \in N_1((A))$ as $n \to \infty$, $[26, Prop.3.2]$ implies that $u_i^n(x) \to u_i(x)$ converges uniformly for $x \in N_1(\Lambda(A))$. Given $\in > 0$, choose $n$ sufficiently large so that $|h_i^n(x) - h_i(x)| \leq \in$ for all $x \in N_1(\Lambda(A))$ and $|\int_c h_i^n(x)\,d\mu(x) - \int_c h_i^n(x)\,d\mu_n(x)| \leq \in$.

Then

$$\lambda_i(\mu) \leq \int_c |h_i(x) - h_i^n(x)|d\mu(x) + |\int_c h_i^n(x)\,d\mu(x) - \int_c h_i^n(x)\,d\mu_n(x)|$$

$$+ \int_c h_i^n(x)\,d\mu_n(x)$$

$$\leq 2\in$$

Since $\in > 0$ is arbitrary, $\lambda_i(\mu) \leq 0$ for all $i$ and $\mu$ is saturated.

Using this lemma, we prove that if K is saturated, then there exists a saturated invariant measure $\mu$ supported by K. Assume K is saturated. Work of Hofbaue and So $[7, Thm.2.1]$ implies that either there exists $y \in W^s(K) \cap C^+$ or K is not isolated for the unrestricted flow $x.t$ If there exists $y \in W^s(K) \cap C^+$, then for all $t \leq 0$ define $v_t = \frac{1}{t}\int_0^t \delta_{y.s}\,ds$ where $\delta_{y.s}$ denotes a Dirac measure based at the point $y.s$. Dissipativeness of (2) and weak* compactness of the Borel probability measures supported by (A) imply there exist $t_k \to \infty$ such that the sequence $v_{tk}$ converges in the weak* topology to a Borel probability measure $\mu$ with support in K.A standard argument implies that $\mu$ is $x.t$ invariant. Define $h_i: C \to R$ by $h_i(x) = (A_i(x)u_i(x), u_i(x))$.Proposition 2 and weak* convergence imply that

$$0 \geq (\lambda i(y) = \limsup_{t\to\infty} \frac{I}{t}\int_0^t h_i(y.s)ds$$

$$= \limsup_{t\to\infty} \int_c h_i\,dv_t$$

$$\geq \limsup_{t\to\infty} \int_p h_i\,d\,v_{tk} = \lambda i(\mu)$$

For all $i$. Alternatively suppose that K is not isolated for the semi flow Then there exists a sequence of positive $\omega - limit$ sets that accumulate on K. Let $\mu_n$ be a sequence of ergodic probability measures supported by these $\omega -$ limit sets. Propositions 1 and 3 imply that $\mu_n$ are saturated for all $n$ Applying Lemma I with $\delta_n$=0 for all $n$ implies that there exists a saturated invariant measure $\mu$ suppoorted by *K*. Hence we have shown that if *K* is saturated then there exists a saturated invariant probability measure $\mu$ supported by *K* Equivalently, (4) holding for all invariant probability measures with support in *K* implies that *K* is unsaturated.

Next we show that if *K* not robustly saturated, then there exists a saturated invariant measure supported by *K* Indeed suppose *K* is not robustly saturated, Then there exists a non-negative sequence $\delta_n$ converging to zero and a sequence of saturated measures $\mu_n$ for $\delta_n$-perturbations of (2)with support in the continuation of *K*. Let $\mu$ be a weak* limit point of $\{\mu_n\}\infty_{n=1}$ Lemma 1 implies that $\mu$ is saturated. Moreover, since the continuation of *K* converges to *K* as $\delta_n \to 0, \mu$ is a saturated invariant measure for (2) supported by *K* Hence we have shown that if *K* is not robustly saturated then there exists a saturated invariant measure supported by *K* Equivalently if (4) holds for all invariant measures with support in *K* then *K* is robustly unsaturated.

To see the equivalence of the conditions given by (4) and (5) let $\Delta = \{p \in R_+^k : \sum_i P_i = 1\}$ and notice that

$$\text{Min max } \lambda i\ (\mu = \min_{\mu} \max_{p \epsilon \Delta} \sum_i Pi\lambda i(\mu)$$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K* The Minimax theorem (see, e. g.,[27] implies that

$$(6)\ \text{min max } \lambda i(\mu) = \min_{\mu} \max_{p \epsilon \Delta} \sum_i Pi\lambda i(\mu)$$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K*. Since min$\mu$

$\sum_i P_i \lambda_i(\mu)$ is attained at an ergodic probability measure with support in *K*, the equivalence of the conditions given by (4) and (5) is established.

To prove the final assertion of the Theorem assume $x \to A(x)$ is twice continuously differentiable and there exists a saturated ergodic measure $\mu$ supported by *K* Proposition 3 implies that there exists *I*C {1,…,k} such that (□$i \in iC_i^+$)=1and $\lambda i(\mu) = 0$ For all $i \in I. SinceKc\,.\partial$□,$\{1, \ldots . k\}\backslash I$ is non-empty. We will show that (2) is not robustly permanent by proving that for all $\delta > 0$ there exists a $\delta$ −perturbation of (2) for which the continuation of *K* is saturated Let $\delta > 0$ be given Choose V $c$ N$\delta$ (*Γ*(A)) to be a compact neighborhood of Γ(A) such that V.t $c$ int *V* for all t > 0.Let $\eta > 0$ be such that *W*=:*N* $\eta$(Γ(A))$c$ *V* Let p:c→ $[0,1] be$ a$C^{\infty}$function such that $p(x) = 1 \quad for\ all\ x \in \Gamma(A) and\ p(x) = 0 for\ all\ x \in C\backslash W\ Define\ \tilde{A} = (\tilde{A}_1, \ldots, \tilde{A}_k)$ by

$$\tilde{A}_i(x) = \{ \frac{A(x)}{A(x) - \frac{\delta}{2Idp\,(x)}} \quad \frac{if\ i \epsilon I}{if\ i \notin I}$$

Where Id denotes the identity matrix of appropriate dimension Let x. t denote the semi flow of $x = \tilde{A}(x)x$.Since$x.t = x.t$ $whenever\ x.[0,t] \in C\backslash W, it\ follows\ that \quad \tilde{V}.t \subset intV\ for\ all\ t >$ $0. Hence, \Gamma(\tilde{A})c\ V\ c\ N\delta\ (\ (A))). We\ also\ have\ ||A(x) - \tilde{A}(x)|| \leq \frac{\delta}{2}$Therefore $x = \tilde{A}(x) is\ a\ \delta -$ perturbation of (2) By construcation,$x.t = x.t\ for\ all\ x \in K and\ t \geq 0. Hence, \mu$ is ergodic for the semiflow of $\dot{x} = \tilde{A}(x)x\ and\ \lambda_i(\mu) \leq -\frac{\delta}{2}$ for this semiflow and $i \in \{1, \ldots, k\}\backslash I$.Let $\mathcal{L}$ and O be the Lyapunov exponents and Oseledec regular points supported by $\mu$ for $x = \tilde{A}(x)x(see, e.g.,$

[28] for definitions).At each point $x \in O$, the splitting of $R^n$ determines three subspaces: the stable subspace $E^s(x)$, the center subspace $E^c(x)$, and the unstable subspace $E^u(x)$ Proposition 3 and our choice of Ã imply that $E^s(x) \cap C^+ \neq \emptyset$. The Pesin stable manifold theorem [28, Corollaries 3.17 and 3.18] implies that tangent to $E^s(x)$ $E^c(x)$ and $E^u(x)$ are locally $x.t$ −invariant families of $C^1$ discs $W^s_x, W^c_x$ and $W^u_x$ corresponding to the stable ,center and unstable manifolds The family of stable manifolds $W^s_x$ is contained in $W^s$ (Ã,supp $(\mu)$). $Since\ E^s(x) \cap C^+ \neq \emptyset$, $W^s_x \cap C^+ \neq \emptyset$ for some $x \in K$. Consequently *K* is saturated for $x = \tilde{A}(x)x$.

The difference between the sufficient and necessary condition is that for the sufficient condition (4) has to hold for all invariant measures supported by *K*, while for the necessary condition (4) has to hold only for ergodic measures supported by *K*. since the invariant measures lie in the convex hull of the ergodic measures the sufficient condition can be more restrictive then the necessary condition.

For three Classes of invariant sets the necessary and sufficient conditions coincide Recall a compact invariant set *K* for (2) is uniquely ergodic if K only supports one invariant measure Recall a compact invariant set *K* is Axiom A if the flow of (2) restricted to *K* is transitive and hyperbolic (see, e.g.,[29] for definitions.)

**Corollary 1:** Assume $x \to A(x)$ is twice continuously differentiable It K$\subset \partial$□ is a compact invariant set and either is uniquely ergodic or Axiom A, then *K* is robustly unsaturated if and only if (4) holds for all ergodic, measures $\mu$ with support in *K*

Proof. If K is uniquely ergodic then the assertion follows immediately For an Axiom A invariant set Sigmund [30, thm.1] has proven that invariant measures supported by periodic orbits of *K* are dense in the set of invariant probability measures supported by *K* In particular any invariant measure can be approximated by an ergodic measure and consequently the result follows.

Another special case where the necessary and sufficient conditions coincide is when *K* supports all populations except one for discrete time models this case was considered by [22].

**Corollary2:** Assume $x \to A(x)$is twice continuously differentiable Let K$c$ $II_i \geq 2C_i^+$ be a compact invariant set. Then K is robustly unsaturated if and only if $\lambda_1(\mu) > 0$ for all invariant measures supported by *K*

## 4. MORES DECOMPOSITIONS AND ROBUST PERMANENCE

To state the sufficient condition for robust permanence, we use a characterization of permanence due to Garay [8] and Hofbauer and So [7] that involves Morse decompositions of the boundary flow Conley[31] defined a collection of sets$\{M_{1,\ldots,}M_m\}$ to be a Morse decomposition for a compact invariant set *K* if $M_1\ldots,M_m$ are pairwise disjoint, compact isolated invariant sent for the flow of (2) restricted to K. For each $x \in K$ there are integers r=r$(x) \leq s =$

$s(x)$ such that $a(x) \subseteq M_r$and $\omega(x) \subseteq M_{s.}$ If $r(x) = s(x), then\ x \in M_{r.}$ Garay, Hofbauer and so [7,8] proved the following characterization of permamnce

**Theorem 2** (Garay, Hofbauer-so) if $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$then (2) is permanent if and only if each of the components $M_1$ are unsaturated.

Theorem 1 and 2 imply the following result: As stated in the above all these results are true for most of problems, interactions, transactions in cosmology, econometrics, finance and other subjects, which we shall illustrate as we go on constructing models for the same.

**Theorem 3:** If $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$and (4) holds for each of the components of the More decomposition, then (2) is robustly permanent Conversely if $x \to A(x) is\ twice\ continuously\ differentiable$

*and (4) is violated by an ergodic measure supported by one of the components of the Morse decomposition, then (2) is not robustly permanent.* Theorem 3 in conjunction with Corollaries 1 and 2 yield a characterization of robust permanence for a class of structured models.

**Corollary 3:** *Let $\{M_1,\ldots..,M_m\}$ os a Morse decomposition for* $\partial$C. *Assume* $x \mapsto A(x)$ *is twise continuously differentiable and for each Morse component $M_i$ one of the following assertions hold $M_i$ is Axiom A, $M_i$ is uniquely ergodic, or There exists* $j \in \{1,\ldots.,k\}$ *such that* $M_i$ ----- $\{x \in C : x_l \gg 0 \text{ for } l \neq j\}$ *Then (2) is robustly permanent if and only if (4) holds for all ergodic measures supported by* $U_i M_i$.

Our result also provided a structured analogue for characterizing totally permanent systems[32]. For an ergodic probability measure$\mu$, define sup ($\mu$) to be the subset $I$ -----$\{1,\ldots,k\}$ such that $\mu(---_{i \in I} C_+) = 1$

**Corollary 4:** *The following statements are equivalent: For all ergodic probability measures $\mu$ with support in* $\partial$C, $\lambda_i(\mu) > 0$ *for all* $I \in \{1,\ldots,k\} \setminus supp(\mu)$ *(2) and all of its subsystems are robustly permanent.*

*Proof. Suppose the first statement holds. Then unsaturated condition given by (5) with p= (1,…,1) and Theorem 3 implies that (2) and all its subsystems are robustly permanent.. The other direction follows form the necessary condition for robust permanence in Theorem 3 and the second assertion of Proposition 3.*

*To illustrate the broad applicability of theorem 3, we develop applications to spatially structured models in section 5 and to a disease model and a gene network model in section 6.*

## 5. APPLICATIONS: PATCH MODELS

A fundamental application of our result are to spatially structured models, with *k* species dispersing between *m* patches {33-35}:

(7) $\frac{dx_i^j}{dt} = x_i^j f_i^j(x^j) + \sum_l d_i^{jl} x_i^l - e_i^j x_i^j$ i = 1,..,k j=1,…,m

*l*

Where $x_i^j$ denotes the density of species I in patch *j*, $x^j = (x^j_1,.....,x^j_k)$ is the vector of species densities in patch *j*, $f^j_i$ is the per-capita growth rate of species *i* in patch *j*, $d_i^{jl} \geq 0$ is the dispersal rate for species *i* from patch *l* into patch *j*, and $e_i^j$ is the emigration rate of species *i* out of patch *j*. Hence $e_i^j \geq \sum_l d_i^{lj}$. We assume that the matrices $(d_i^{lj})_{j,l}$ are irreducible for each species *i* . it is then easy to write (7) in the form (2), with $x_i = (x^l_i, ..., x_i^m)$ 'where' denotes transpose.

For a single species (7) generates a monotone flow. Under mild assumption, e.g., each $f^j_i$ *is* decreasing and negative for large densities, there is a globally stable equilibrium, see, e.g.,[36] and [37, sect. 5.4]. This equilibrium will be positive if the invasion rate at the origin λ (0), which is given by the leading eigenvalue (stability modulus) of the matrix $(d^{jl} + f^j(0) - e^j)$, is positive.

5.1 Two types of debitable and creditable systems:. For *two competing species*, e.g., $f^j_i$ is decreasing with respect to both species and is negative at large densities, the dynamics are still monotone, and hence almost all orbits converge to an equilibrium. Robust permanence requires two conditions. First, $\lambda_i$ (0) >0 for both species in which case there are two single species equilibria $E_1$ and $E_2$. Second, the invasion rates $\lambda_2(E_1)$ and $\lambda_2(E_2)$ are positive. Whether there is a unique positive globally stable equilibrium for the robustly permanent system depends in a delicate way on the

system parameters, see [38]. For slowly dispersing populations (i.e. $0< e^j_i \ll 1$ for all *i, j* ), the robust permanence condition is particularly

Straight forward to verify. Let $E^j_i$ be the largest solution to $x^j_i \quad f_i{}^j \ (x^j_i) = 0$ i.e. the equilibrium attained by species *I* in patch *j* when there is no dispersal and no competitors. For slowly dispersing populations, $E_1$ is close to the product $(E_1{}^1, \ldots, E_1{}^m)$, and the matrix $A^2(E_1)$ is close to the diagonal matrix $\text{diag}(f^i{}_2(E_1{}^j))^m_{j=1.}$ Hence the invasion rate $\lambda_2(E_1)$ is close to $\max_j f_2^j(E^j_1)>0$. So we obtain: The system with two slowly dispersing competing species is robustly permanent if and only if system refers to system of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

(8) $\max_j f_2^j(E^j_1)>0$ and $\max_j f_2^j(E^j_2)>0$

Intuitively, robust permanence requires for each species there is at least one patch where it can persist.

For *systems, particle interactions, financial transactions, cosmological interactions and reactions, as we shall show model by model the* situation are even easier. Since (under mild assumption on the prey dynamics) the Morse decomposition of $\partial C$ is given by two equilibria $\{0,E_1\}$, the spatial predator-prey system is robustly permanent, if the predator can invade at the prey equilibrium: $\lambda_2(E_1) > 0$. However, the global dynamic are likely to be more complicated. For instance, if in each patch there is a globally stable limit circle (e.g.,[39-41]) and dispersal is sufficiently slow (i.e. $0< e^j{}_i \ll 1$for all *i, j*), our result about robust permanence and the theory of normally hyperbolic manifolds [42] imply that there is a positive *m* dimensional torus, which attracts almost all orbits in C, whenever $\max_j f_2^j(E^j_1)>0$

**5.2 Rock-Paper-Scissors Dynamics.** The Lotka-Volterra model of rock-paper-scissor dynamics is a simple model that is used as property for understanding intransitive ecological outcomes [43,44]. Here, a simple spatial version of this dynamic is given by

$$\frac{dx_1^j}{dt} = x_1^j(1 - x_1^j - \beta^j x_2^j - \alpha^j x_3^j) + \sum d_1^{jk} x_1^k - e_1^j x_1^j$$

$\frac{dx_2^j}{dt} = x_2^j(1-\alpha^j x_1^j - x_2^j - \beta^j x_3^j) + \sum d_2^{jk} x_2^k - e_2^j x_2^j$

$\frac{dx_3^j}{dt} = x_i^j(1-\beta^j x_1^j - \alpha^j x_2^j - \beta^j x_3^j) + \sum d_3^{jk} x_3^k - e_3^j x_3^j$

Where $a^j \in (0,1)$ and $\beta^j > 1$ for all $j$. A more general version of this dynamic is presented in [44]. Under the assumption that there is no cost to dispersal (i.e. $e^j_i = \sum_l d^{l,j}_i$ for all $i, j$) , the maximal invariant set in $\partial C$ consist of the origin 0 and a heteroclinic cycle connecting positive single species equilibria$E_1$, $E_2$ and $E_3$. For slowly dispersing population ($0 < e^j_i \ll 1$ for all $i$, j), $E_1$ is close to $(x_1, x_2, x_3) = (1,\ldots,1,0,\ldots,0,0,\ldots,0)$ and the invasion rates $\lambda_2$ ($E_1$)and $\lambda_3$ ($E_1$)are close to $\max_j 1-a^j$ and $\max_j 1-\beta^j$. For the other equilibria $E_i$, the debitable rates of the missing species are also given by $\max_j 1-a^j$ and $\max_j 1-\beta^j$. Consider the Morse decomposition of $\partial C$ given by 0 and the heteroclinic cycle. Since $\lambda i\ (0) = 1 > 0$ for $i = 1, 2, 3, 0$ is robustly insaturated. An algebraic computation reveals that (4) holds for all invariant supported by the heteroclinic cycle (i.e. all convex combination of the Dirac measures supported by the equilibria) if and only if

$$\max_j(1-a^j) > -\max_j(1-\beta^j)$$

Equivalently,

$$2 > \min_j a^j + \min_j \beta^j$$

In particular, even if the heteroclinic cycle is attracting for each patch when the system is uncoupled (i.e. $2 < a^j + \beta^j$ for all $j$ ), it can be repelling for the weakly coupled system.

5.3. Three species **Lotka Volterra in spatially homogenous environments.** Consider Lotka-Volterra dynamics in a spatially homogenous environment: $(x+a)^n =$

(9) $$\frac{dx_i^j}{dt} = x_i^j\left(r_i - \sum_s a_{is} x_s^j\right) + \sum_l d_i^{jl} x_i^l - e_i^j x_i^j, \qquad i = 1,\ldots,k;\ j = 1,\ldots,m$$

We assume that dissipative can be shown by a linear Liapunov function: There are $c_i > 0$ such that $\sum_i c_i\, a_{is} x_i x_s < 0$ holds for all $x \in R_+^k \setminus \{0\}$. Then the weighted sum of densities across all patches $\sum_{i,j} c_i\, x_i^j$ is decreasing for large densitities and hence (9) generates a dissipative semiflow with global attractor $r$. We further assume that for each patch *j* and each species *i* the immigration rate equals the emigration rate , i. e,

(10) $$e_i^j = \sum_l d_i^{jl}$$

This guarantees that the set $H = \{x \in C: x_i^j = x_i^l \forall i, j, l\}$ of spatially homogenous states is forward invariant under (9).

Now consider the Lotka-Volterra dynamics without spatial structure

(11) $$\frac{dx_i}{dt} = x_i(r_i - \sum_s a_{is} x_s), \qquad i = 1, \ldots, k$$

(11) $$\frac{dx_i}{dt} = x_i(r_i - a_{is} x_s), \quad i = 1, \ldots, k$$

It has been shown [12,ch.16.1,16.2] That for *k* =3,(11) is robustly permanent (robust meaning here within the class of Lotka –Volterra systems) if and only if all equilibria on the boundary of $R_+^3$ are unsaturated and whenever there is a heteroclinic cycle connecting the one species equilibria then this cycle is repelling. Equivalently the following four conditions hold. There exists an interior equilibrium $\dot{x}$ (i.e $A\dot{x} = r \;\; with\; \dot{x} \gg 0$); Det (-A)>0; the 2 species subsystems are not bistable competition systems if there is a heteroclinic cycle between the one species equilibria, say $\mathbf{E_1 \to E_2 \to E_3 \to E_1}$ then the following inequality holds

12) $$\lambda_2(E_1)\lambda_3(E_2)\lambda_1(E_3) > |\lambda_3(E_1)\lambda_1(E_2)\lambda_2(E_3)|$$

Furthermore, the proof in [12, ch.16.1, 16.2] shows that the boundary flow of (11) has a simple Morse decomposition: if (iv) applies then there are two Morse sets, the origin 0 is a repeller, and the heteroclinic cycle is

the dual attractor (within$\partial R^3_+$); if (iv) does not apply then the finitely many boundary equilibria form a Morse decomposition of the boundary flow.

**Species, population we mean types of and population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.**

**Theorem 4:** Under the above assumptions,(9) for k=3 is permanent if and only if *(i),(ii),(iii)* and *(iv)* hold.

Proof: Suppose (9) is robustly permanent. Since the nonspatial system (11) is the restriction of (9) to the invariant subspace **H**, (11) must be robustly permanent and hence*(i)-(iv)* hold

Conversely, suppose that (i)-(iv) hold. We show that the above Morse decomposition of the boundary flow of (11) (now in **H**) is also Morse decomposition for (9) restricted to the maximal invariant subset of $\partial C$ which is contained in the spatial two species subsystems of (9). For this we show that in each 2 species system the orbits of (9) converge to **H** (For the one species subsystems this is obvious.) If a two species subsystem of (11) has an internal equilibrium, say $\mathbf{E_{12}}$ then by assumption (ii),this is not a saddle but a sink, and hence $a_{11}a_{22} > a_{12}a_{21}$. This implies that the 1-2 submatrix of A is VL-stable [12,ch.15.3,],and [45] implies that $\mathbf{E_{12}}$ is the global attractor for the spatially structured system (9) restricted to the first two species.

If a two species subsystem of (11) has no internal equilibrium, then a one species equilibrium say $\mathbf{E_1}$ is the global attractor. Then Lemma 2 (see below) show the corresponding result for (9).A similar proof applies to the case where the origin **0** is the global attractor of a two species subsystem.

Finally observe that under the above assumptions (9)-(10), the invasion rates $\lambda_i$ at spatially homogenous boundary equilibria are the same for (9) and (11) Since all boundary equilibria are unsaturated for (11), they are for (9) as well. The same applies in case (iv) to the heteroclinic cycle. Note that the condition (12) on the eigenvalues is equivalent to the fact that all invariant measures supported on the heteroclinic cycle are unsaturated, see [16,Ex.4.5]

Species we mean types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

**Lemma 2:** Suppose, in a two species Lotka –Volterra system (11) (k=2) all interior orbits converge to the one species equilibrium $\mathbf{E_1}$ .Then all orbits of the spatial version (9) (k=2) converge to the spatially homogeneous one species equibrium $\mathbf{E_1} \in \mathbf{H}$.

**Proof:** If the system is of predator-prey or mutualistic type then Hasting's [45] result applies again. So let us assume that the local interaction is given by the competition system

(13) $\frac{dx_1}{dt} = r_1 x_1(1 - x_1 - ax_2)$ $\frac{dx_2}{dt} = r_2 x_2(1 - \beta x_1 - x_2)$

With 0<a <1< $beta$.In this case, Kazuo Kishimoto (letter to JH, Oct 1987) has given the following argument. There exists a family of forward invariant rectangles contracting to $\mathbf{E_1}$. This shows that all interior solutions of (9) (k=2) converge to $\mathbf{E_1}$.

Theorem 4 shows that a spatial network of identical patches is permanent, if the within patch dynamics is permanent, and three species Lotka-Volterra dynamics. It is not clear how to extend this result to more than three species. The crucial step in the three species case is that permanence precludes bistable two species subsystems. The spatial version of a bistable two species system allows plenty of stable equilibria outside **H,** see [33].However, permanent four species Lotka Volterra systems may have two species bistable subsystems [12,ch.16.4].So there may be invariant sets on the boundary outside **H** which need to be unsaturated to make (9) permanent.

## 6. MORE APPLICATIONS

6.1. **Disease, particle, product, marketing, financial, econometric dynamics with density-dependent demography**. GAO and Hethcote [46] introduced a model of disease dynamics with density-dependent demography. Here we consider a variation of their model in which a population of size P has a per-capita birth rate b(P) and per –capita death rate d. To ensure that the population persists in the absence of the disease, we assume that there exists K>0 such that b(K)=d. To describe the disease dynamics, let *S* be the number of individuals susceptible to the disease, *I* the Number of infected individuals, and R the number of individuals that have recovered from the disease. We assume that P=S+*I* +R. On the same lines we can think of sick and non performance accounts, and those which are conducting themselves well, and those which are satisfactory and find a parallel formula. The same thing holds true in case of a black hole eating matter and black hole radiating radiation.. If the disease transmission is asymptotic, β is the contact rate between susceptible and infected individuals, γ is the rate at which individuals recover from the disease, and m the mortality rate due to disease, then the population dynamics is given by

$$\frac{dp}{dt} = b(P)P - dP - mI$$

$$\frac{dI}{dt} = \beta(P - I - R)\frac{I}{\epsilon + P} - (\gamma + d + m)I$$

$$\frac{dR}{dt} = \gamma\, I - dR$$

Where $\epsilon>0$ is a constant It is useful to introduce a change of coordinates in which $y=\frac{I}{P}$ and $z=\frac{R}{P}$:

$$\frac{dP}{dt} = (b(P)-d-my)P$$

$$\frac{dy}{dt} = \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(P)\ y$$

$$\frac{dz}{dt} = \gamma y - (b(P) - my)z$$

Setting $x_1$=P and $x_2$=(y,z) 'where' denotes transpose yields a structured model where $A_1(x)$=b-d-my and

$$A_2(x) = \begin{pmatrix} \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(p) & 0 \\ \gamma & my-b(p) \end{pmatrix}$$

Since {0,(*K*,0,0)}is a Morse decomposition of the boundary dynamics Theorem 3 and Remark 1 imply that this model is robustly permanent if and only if b(0)>d and $\beta\frac{K}{\epsilon+K} > \gamma + m + d$. In particular, for $\epsilon > 0$ sufficiently small, one requires that the basic reproductive number $\frac{\beta}{r+m+d}$ is greater then one.

Remarkably this same criterion determines robust permanence of model with significantly more complicated boundary dynamics. For example, suppose the focal population is a predator species. If the prey has abundance *N* and exhibits logistic dynamics f(*N*)is the per-capita predator consumption rate of the prey and b(N)is the per-capita reproductive rate of the predator then the dynamics become

$$\frac{dN}{dt} = rN(1 - N/K - f(N)P$$

$$\frac{dP}{dt} = (b(N) - d - my)P$$

$$\frac{dy}{dt} = \beta(1 - y - z)\frac{P}{\epsilon+P} - r - m(1-y) - b(N)\ y$$

$$\frac{dz}{dt} = \gamma y - (b(N) - my)z$$

Where r>0 and *K*>0, and *f*(*N*),b(N)satisfy f(0)=b(0)Define $x_1$=*N*, $x_2$=*p*, $x_3$=*(y,z)*′ Let us assume that the disease-free community (i.e.y=z=0)is permanent (i.e.b(K)>d) and let A be the global attractor in the interior of the P-N plane Under these assumption a Morse decomposition of the boundary dynamics is given by $(x_1, x_2, x_3) = (0,0,0)$, $(x_1, x_2, x_3) = (K, 0,0)$and A. Since the equilibria (0,0,0) and (k,0,0) are unsaturated, it remains to characterize $\curlywedge_3(\mu)$for any invariant measure μ supported by A. Proposition 3 and Remark 1 imply

$$\curlywedge_2(\mu) = \int_c (bx_1)\text{-}d)d\mu(x) = 0$$

*Hence* $= \int_c bx_1)d\mu(x) = d$ *and*

$$\curlywedge_3(\mu) = \beta \int \frac{x_2}{\epsilon + x_2} d\mu(x) - r - m - d$$

**Theorem 3**: and Remark 1 implies that this structured model is robustly permanent if $\frac{\beta}{r+m+d} > 1 and\ \epsilon > 0$ is sufficiently small.

6.2. Gene networks. The repressilator is an oscillatory gene network based on three (or more generally an odd number of transcriptional repressors. Such a system has been genetically engineered in E-coli by Elowitz and Leibler [47] Mathematical model go back to [48.49] they all involve concentrations of proteins and mRNs. A modified model with auto-activation was suggested in [50]

$$(14)\ x_i = \beta(y_i - x_i)$$

$$(15)\ y_i = aF(x_{)i} \text{-} y_i)$$

With

$$(16)\quad F(x)_i = x_i g\ (x_i, x_i - 1) = \frac{x_i}{1 + x_i + p x_{i-1} + K p\, x_i x_i - 1}$$

Here $x_i$ and $y_i$ are normalized concentrations of proteins and mRN As belonging to gene $i$,p is the strength of repression and k a cooperativity parameter this is a structured system of type(2)

$$(17)\ \begin{pmatrix} x_i \\ y_i \end{pmatrix} dot = \begin{pmatrix} -\beta & \beta \\ ag(x_i,x_i)-1 & \end{pmatrix} \begin{pmatrix} x_i \\ y_i \end{pmatrix}$$

In [50,Thm 3] it was shown that for n=3,a>1 and p>1,system(14)has a heteroclinic cycle connecting the 3 single gene equilibria (similar to the rock –scissors-paper dynamics of section5.2)Moreover

- the system is permanent i.e. the boundary of $R^6$+is repelling, if $\lambda + \mu > 0$
- this heteroclinic cycle is asymptotically stable if $\lambda + \mu < 0$

where $\lambda > 0$ and $\mu < 0$ are the invasion rates of the two missing genes i.e. the leading eignvalues of the 2X2 matrix in (17)evaluated at the appropriate single gene equilibria (for explicit expressions see[50,eqs.(127,128)].)

Theorem 3 implies the permanence condition implies the stronger conclusion of robust permanence A similar result holds for n odd, but the heteroclinic cycle is then actually a heteroclinic network.

Note: For references please see the original paper.

One of the examples for which we apply the model are as follows: **Progress in Particle and Nuclear Physics Volume 88, May 2016, Pages 1–28 Chiral magnetic and vortical effects in high-energy nuclear collisions—A status report D.E. Kharzeeva, b, J. Liaoc, d, , , S.A. Voloshine, G. Wangf**

The interplay of quantum anomalies with magnetic field and vorticity results in a variety of novel non-dissipative transport phenomena in systems with chiral fermions, including the quark–gluon plasma. Among them is the Chiral Magnetic Effect (CME)—the generation of electric current along an external magnetic field induced by chirality imbalance. Because the chirality imbalance is related to the global topology of gauge fields, the CME current is topologically protected and hence non-dissipative even in the presence of strong interactions. As a result, the CME and related quantum phenomena affect the hydrodynamical and transport behavior of strongly coupled quark–gluon plasma, and can be studied in relativistic heavy ion collisions where strong magnetic fields are created by the colliding ions. Evidence for the CME and related phenomena has been reported by the STAR Collaboration at Relativistic Heavy Ion Collider at BNL, and by the ALICE Collaboration at the Large Hadron Collider at CERN. The goal of the present review is to provide an elementary introduction into the physics of anomalous chiral effects, to describe the current status of experimental studies in heavy ion physics, and to outline the future work, both in experiment and theory, needed to eliminate the existing uncertainties in the interpretation of the data. Keywords Chiral magnetic effect; Chiral vortical effect; Chiral anomaly; Quark-gluon plasma; Heavy ion collisions

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# Sufficient Conditions For The Existence And Linear Stability Of Time Independent Age Distributions In Respect Of Particles, Propositions, Problems, Hegelian Structures, Deleuzean Treatises, Interactions And Transactions In Cosmos And Others

**Abstract:** Gopalaswamy studied sufficient conditions are obtained for the existence and linear stability of time independent age distributions in two species competition with age time lagged density dependent mortality and fertility functions.. Work was published in Bulletin Australian Mathematical Society Volume 25 (1982). We show how the model could be applied in case of finance, econometrics, psycho economics, interactions between particles and transactions in nature. Basically the competition is for resources, notwithstanding such experiments could be conducted in laboratory like a rat kept in front of a cat, such incidents could be represented mathematically, and the corresponding linear stability conditions ca be obtained in case of interactions in cosmos. Take for example the case of a cash credit account of a borrower. Many inward bills waiting for paid, or inland or foreign bills payment all are in line for the debit of the account, which also acts as a resource. In addition there are recurrent cases where in the credit is issued to the cash credit account in case of demand and inland bills. In case of economics, investment has inventories and it has to be utilized towards the end of consummation of output. And this is case of output gobbling part of investment which is inventories. Output has to be sold to the customer and the amount received is diverted partly to savings and consumption and part could be employed towards the consolidation of the system. So this is a case of investment itself acting as a both receiver and payer. It shall eat and also produce the finance. Savings as stated above could be deployed for the further diversification and proliferation and that is towards the end of putting the company on terra firma. Another example in cosmology is the case of black holes guzzling matter. There are lot of black holes and lot of matter. Result of such a guzzling the production of Hawking radiation. Horava gravity is another example. How the conformal field theory could be changed in the case of Horava gravity. Still another case is the case of Vonneumann algebra in the cae of Horava gravity. In this case like cash credit account, the Horava gravity and Von Neumann algebra act as both receivers and takers. Such incidents are legion in cosmology and even in philosophy. We have considered self as individual general ledger, and the debits and credits are raised to the same account thus anointing as both a receiver and payer. We give plethora of such cases and apply the model. Due to spatial constants, same module is used for all the modules and we rewrite to state that the model is applicable to all modules. This is due to spatial restrictions. What is to be done is that each and every parameter is to be changed with a subscript or a superscript and finally it is to be concatenated towards the end of obtention of prediction model or optimizing model, for Duality of motion model for further analysis. Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a

qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

**Introduction**

Let $P_1$ *(t)* and $P_2$ *(t)* denote the total particles, interactions, transactions, etc, which we dilate upon further investigation. Perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. **Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems** sizes (or biomasses) at time $t\geq 0$ of two interacting species living in a common habitat and competing for a common pool of resources. The competition for resources will be implicit in our model similar to that in the two species Lotka- Volterra competition system. Assuming constant sex ratios in the two species we can consider $P_1$ *(t)* and $P_2$ *(t)*to be the Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

population of f; immigration, emigration and internal dispersion in the habitat are assumed to play no significant role in the dynamics of the community.Popoulation we mean Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

We suppose that the two species contain respectively $p_1(a, t)\,da$ and $p_2(a, t)\,da$ individuals with ages between $a$ and $a + da\ (a\geq 0)$ at time $t$ so that we have

(1.13) $P_1(t) = \int_0^\infty \rho_1(a, t)\,da,\ P_2(t) = \int_0^\infty \rho_2(a, t)\,da,\ t\geq 0.$

The rates of change of the two species densities are by definition $D\rho_1(a, t)$ and $D\rho_2(a, t)$ where

(1.14) $D\rho_i(a, t) = \limsup_{h\to 0} \frac{\rho_i(a+h,t+h)-\rho_i(a,t)}{h},\ i = 1, 2.$

Assuming the existence of the limits in (1.2) we consider the following time lagged model system

(1.15) $$\begin{cases} D\rho_1(a,t) + f_1(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_1(a, t) = 0, & a>0,\ t>0 \\ D\rho_2(a,t) + f_2(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_2(a, t) = 0,' \quad (1.4) \\ \rho_1(0, t) = \int_0^\infty b_1(a, P_1(t), P_2(t-\tau))\rho_2(a, t)da, & t>0, \\ \rho_2(0, t) = \int_0^\infty b_2(a, P_1(t), P_2(t-\tau))\rho_1(a, t)da, \end{cases}$$

(1.5) $\rho_i(a, s) = \varphi_i(a, s),\ a.0,\ s\epsilon[-\tau, 0],\quad i = 1, 2,$

(1.6) $P_i(s) = \int_0^\infty \varphi_i(a, s)\,da,\ s\in[-\tau, 0],$

Where $\tau$some is fixed nonnegative constant; $f_1$ and $f_2$ denote the age and density dependent mortality rates with time lags in interspecific interactions, $b_1$ and $b_2$ denote the age and density dependent fertility functions again having time lags in the interspecific interactions; $\varphi_1$ and $\varphi_2$ denote the initial age distributions needed for the formulation of the model. Note we mean the population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

The model system (1.1)-(1.6) will be meaningful only if $\varphi_1$, $\varphi_2$, $f_1$, $f_2$, $b_1$, $b2$ are nonnegative functions of the respective arguments. Also since the total initial populations have o be finite, $\varphi_1(\cdot, s)$ should belong to $L_1(R^+)$ for each $s\epsilon$ $[-\tau, 0]$. That $f_i$ and $b_i$ $(i = 1, 2)$ are independent of t explicitly and this assumption can be interpreted to represent the temporal uniformity of the environment.

One of the fundamental questions for (1.1)-(1.6) is the following; under what conditions on $f_i$, $b_i$ and $\tau$ does the system (1.1)-(1.4) have time independent nontrivial solutions, and when such solutions, and when such solutions exist, are they stable with respect to some suitable stability criterion? When such stable time invariant solutions exist for (1.1)-(1.4) we say that the two types has a stable age distribution. Existence of a stable age distribution is the analogue of the existence of a stable steady state for the corresponding age independent system. Types we mean types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

In this article we investigate the above question for the case of a competitive interaction with some additional constraints on the vital rates $f_i$ and $b_i$ $(i = 1, 2)$. Age dependent population systems without time delays have been considered by several authors (Gurt in and MacCamy [3, 4, 5], Gurt in and Levine [6], Haimovici [7] and Rotenberg [8]). The question of the relation between an age independent system and a corresponding age dependent system has been considered by Gurt in and MacCamy [4] and Gurt in and Levine [6] who have established an asymptotic (*as* $t\rightarrow\infty$ ) relation between such models by constructing a higher dimensional 1 umped parameter system (in terms of ordinary differential equations) to represent the age dependent distributed parameter system. Assuming the existence of stationary age distributions Haimovici [7] considers their stability in a system of two interacting populations explicitly taking into consideration the dynamical nature of the encironments resources and pollution.

In this article we add another realistic feature namely time lags in the age dependent models and obtain sufficient conditions for the existence of stationary age distributions in (1.1)-(1.4) and show that the time lags in interspecific

interactions have no effect on the linear stability of the age distributions although the decay rates of perturbations will depend on the magnitudes of the time delays. Such "harmless" natures of time lags in interspecific interactions have been noted by the author for the age independent models in [1, 2].

## 5. Existence and uniqueness of solutions

We first narrate our assumptions on the vital rates of $f_i$ and $b_i$ $(i = 1, 2)$ so that the system (1.1)-(1.6) will denote a competition system. Let $BC\ (0, \infty)$ denote the linear space of bounded continuous real functions on $(0, \infty)$ with the norm defined by

$$||f|| = \{max\ |\ f(t)\ |,\ t\ \in (0, \infty)\}.$$

Let $K_1$ be the cone of nonnegative functions of $BC\ [0, \infty]$. Let $C\ [-\tau, 0]$ denote the space of real continuous functions with a norm defined by $||g|| = \{\ max\ |g(s)\ |,\ s \in [-\tau, 0]\}$; let $K_2$ be the nonnegative cone of functions in $C\ [-\tau, 0]$ . We present our first set of assumptions on $f_i$, $b_i$ and $\varphi_i$ for $i = 1, 2$.

$A_1$. $f_1 \in C\ (R^+ \times K_1 \times K_2, R^+)$, $f_1 \geq d_1^*$ on $R^+ \times K_1 \times K_2$

$f_2 \in C\ (R^+ \times K_2 \times K_1, R^+)$, $f_2 \geq d_2^*$ on $R^+ \times K_2 \times K_1$ ,

( $d_1^*$ , $d_2^*$ are positive constants);

$f_1$ and $f_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ ; that is, *for* $i = 1, 2$,

$$|f_i\ (a, P_1^{'}\ , P_2^{'}) - f_i(a, P_1^{"}\ , P_2^{"})| \leq d_i\ [||P_1^{'} - P_1^{"}, P_2^{'} - P_2^{"}||]$$

($d_1$, $d_2$ are positive constants),

$$\partial f_1/\ \partial P_1, \partial f_1/\ \partial P_2 \in C(R^+ \times K_1 \times K_2, R^+)\ ,$$

$$\partial f_2/\ \partial P_1, \partial f_2/\ \partial P_2 \in C(R^+ \times K_2 \times K_1, R^+),$$

$A_2$. $b_1 \in C\ (R^+ K_1 \times\ K_2, R^+)$, $b_1 \leq \beta_1^* on\ R^+ \times K_1 \times K_2$ ,

$b_2 \in C\ (R^+ K_2 \times\ K_1, R^+)$, $b_2 \leq \beta_2^* on\ R^+ \times K_2 \times K_1$ ,

*($\beta_1^*$ , $\beta_2^*$ are positive constants);*

$b_1$ , $b_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ *; that is for* $i = 1, 2$,

$$|b_i\ (a, P_1^{'}, P_2^{'}) - b_i\ (a, P_1^{"}, P_2^{"})| \leq \beta_i\ [||P_1^{'} - P_1^{"}|| + ||P_2^{'} - P_2^{"}]$$

$B_1$ , $B_2$ Are positive constants),

$b_i(\circ,(P_1,P_2)\epsilon L^{\infty}(R^+,R^+)$ For fixed $P_1,P_2,i=1,2,$

$$\frac{\partial b_1}{\partial P_2},\frac{\partial b_1}{\partial P_2}\epsilon\ C\ R^{+}\text{x}\ k_{\ 1}\text{x}\ k_{\ 2},\text{R}^{-}),$$

$$\frac{\partial b_2}{\partial P_1},\frac{\partial b_2}{\partial P_2}\epsilon\ C\ R^{+}\text{x}\ k_{\ 2}\text{x}\ k_{\ 1},\text{R}^{-}),$$

$A_3$. $\varphi_i\ (\cdot\ ,\text{s})\in \text{L}^{\ 2}\ (\text{R}^{+},\text{R}^{+})\ \cap \text{C}(\text{R}^{+},\text{R}^{+})$ for each $s\in\ [-\tau,0]$,

$i,\ =1,\ 2,$

By a solution of $(1.1)-(1.6)$ on $[0,\text{T}]$, we mean

$\text{P}_i$: $\text{R}^{+}$ ×$[-\tau,\text{T}]\rightarrow \text{R}^{+}$ $(i,\ =1,\ 2,)$ With the following properties:

$S_1$. $\text{P}_i(\cdot\ ,t)\in \text{L}^{\ 1}\ (\text{R}^{+},\text{R}^{+})$ for each $t\ \in[-\tau,\text{T}]$;

$S_2$. $\text{P}_i(a,\cdot\ )\in \text{C}\ ([-\tau,\text{T}],\text{R}^{+})$ for each $a\ \in \text{R}^{+}$;

$S_3$. $\text{DP}_i$ Exists along the characteristics $t-a=$ constant

On $\text{R}^{+}$ ×$\{\text{R}^{+}\cap\ [0,\text{T}]\}$ And is continuous $(i,\ =1,\ 2,)$;

$S_4$. $\text{P}_i$ Satisfies the system (1.1)- (1.6) for

$(a,t,)\ \in \text{R}^{+}$ ×$[-\tau,\text{T}]$ .

Let us first convert the system (1.1)- (1.6) into an equivalent system of integral equations. For brevity we denote by P the pair ( $\text{P}_1,\text{P}_2$) If we let

$$\text{P}_i(a+x,t+x)=\ \ \text{P}_i(x),(i,\ =1,\ 2,)\ ,$$

Then (1.3) considered at $(1.3)$ considered at $(a+x,\ t+x)$ becomes

$$\frac{dp_{\ i}(\text{x})}{dx}+f_i(a+x,\ P(t+x))\ \ P_i(x)=0$$

Which has a solution of the form?

$$P_i(x)=P_i(0)\ \exp\left[-\int_0^{\infty} f_{\ i}\big(a+s,P(t+S)\big)ds\right]$$

And hence

(2.1) $p_{\ 1}\ (a+x,t+x)=p_{\ 1}\ (a,t)\ exp\big[-\int_o^x\ f_{\ 1}\ \big(a+s,p(t+s)\big)ds\big]$

For $a \geq t$ *we choose (a,t)=(a-t,0)* in (1.3) and *x=t* in (2.1) so that

(2.2) $p_1(a,t) = \varphi_1(a-t,0)\exp\left[-\int_o^t f_1\left(a-t+s, p_1(s)\right)ds\right]$

$$O \leq t \leq a, i = 1,2,$$

With

$$P_1(s)=(p_1(s),p_2(s-t)),$$

$$P_2(s)=(p_1(s-t),p_2(s)),$$

For $0 \leq a < t$ we choose *(a,t)=(0,t-a)in (1.3)and x=a in (2.1)* so that

*(2.3)* $p_i(a,t)= B_i(t-a)exp[-\int_o^a f_1(s, p_i(t-a+s))ds]$

$$t > a \geq, i = 1,2,$$

Where

*(2.4)* $B_i(t) = p_i(0, t), i = 1,2, t > 0.$

Define $M_i and\ L_i\ (i = 1,2)$ as follows:

$$M_i(a,t, p_i(t)) = exp[-\int_o^{t-a} f_1(s, p_i(a-a+s))ds]$$

*(2.5)* $L_i\left(a,t,\ p_i(t)\right) = \exp[-\int_o^t f_1(\text{a+s},\ p_i(s))ds]$

It will now follow from (1.1),(1.3),(2.2)-2.5)that

*(2.6)* $p_i(t) = \int_o^t B_1(a) M_1(a, t, p_i(t)da + \int_o^\infty \varphi_i(a, 0)\ L_i(a,t,p_i(t)da,$

$$t > 0, i = 1,2,$$

*(2.7)* $p_i(s) = \Phi_i(s) = \int_o^\infty \varphi_i(a, s)da, s \in [-\tau, 0], i = 1,2,$

(2.8) $\beta_i(\text{t}) = \int_0^t \beta_i(\text{a}) b_i(\text{t-a}, P_i(t)) M_i.\left(a,t,P_i(t)\right)da$

$$+\int_0^\infty b_i\left(a+t, P_i(t)\right)\varphi_i(a,0)L_i\left(a,t,P_i(t)\right)da$$

$$i = 1, 2, \text{t} > 0$$

the equivalence of (2.6)-(2.8) with (1.1)-(1.6) is established by the following whose proof is identical to a similar result of Gurtin and MacCamy [4]; hence we omit the details of the proof of the following.

THEOREM 1 : *Suppose the system* (1.1)-(1.6) *satisfies the conditions* $A_1, A_2, A_3$ . *If* $(P_1, P_2)$ *is a solution of*

*(1.1)-(1.6) then the total population* $P_1, P_2$ *and the birth rates* $B_1, B_2$ *satisfy the integral equation (2.6)-(2.8).*

*conversely if* $P_1, P_2, B_1, B_2$ *are nonnegative continuous solution of (2.6)-(2.8) on [ 0, T] and if* $P_1, P_2$ *are*

*defined by (2.2)-(2.4) on* $R^+$ x [0, T] *then such* $P_1, P_2$ *provide a solution of (1.1)-(1.6) on* $R^+$ x [0, T] .

The following a *priori* estimates are useful to prove our existence theorem below. It will immediately follow from the bounds for $f_i$ and $b_i$ that

$$M_i(a,t,P) \leq \exp[-d_i^*(t-a)] \leq 1, if \quad t \geq a,$$

*(2.9)*

$$L_i(a,t,P) \leq \exp[-d_i^* t]$$

From (2.7)-(2.8) we derive that

$$\beta_i(t) \leq \int_o^t \beta_i^* \beta_i(a) \exp[-d_i^*(t-a)]\, da + \beta_i^* \Phi_i(0) \exp[-d_i^* t], \quad i = 1,2,$$

And hence by Gronwall's inequality,

(2.10) $$\beta_i(t) \leq \beta_i^* \Phi_i(0) \exp[\delta_i t], \ \delta_i = \beta_i^* - d_i^*$$

(2.5), (2.6) and (2.10) lead to

(2.11) $$|P_i(t) - 0_i(0)| \leq \beta(0)\, exp[\delta_i\, a - d_i^*(t-a)]\, da$$

$$+\int_0^\infty \varphi_i\,(a,0)|1 - exp(-d_i^* t)|da$$

$$\leq \beta_i^* \Phi_i\,(0)\, exp(\beta_i^* t)\left\{\frac{1-exp[-(d_i^*+\beta_i^*)t]}{d_i^*+\beta_i^*}\right\} + \phi(0) d_i^* t$$

$$\leq [\beta_i^* \phi(0)\, exp(\beta_i^* t)\, \phi(0) d_i^*]t \quad for\ t > 0.$$

We can now prove the following result on the existence and uniqueness of solution of the integral equations (2.)-(2.8)

**THEOREM 2:** *Assume that the system* (1.1)-(1.6) *satisfies the hypotheses*

| |
|---|
| $A_1, A_2, A_3$. *Then there exists a unique set of continuous function* $(P_1, P_2, B_1, B_2)$ such *that* |
| $P_i : R^+ \to R^+, \qquad B_i : R^+ \to R^+ \ (i = 1,2\ ),$ |
| *Which satisfy the integral equations* (2.6)-(2.8) |
| Proof. We will first show the existence of local solution on same (possibly small) interval [0, T] $\subset R^+$ such that |
| $P_i \in \mathrm{C}([0,T], R^+), \quad B_i \in \mathrm{C}([0,T], R^+)\ , i = 1, 2,$ |
| Satisfying (2.6)-(2.8) and then show the $P_i$ and $B_i$ can be continued as solution for any finite interval in $R^+$ for some fixed |
| Positive number $T$ we let |
| $C^+[-\tau, T] = \{f \in \mathrm{C}([-\tau, T], R)\,|f \geq 0\}$ |
| Let $\|.\|_T$ denote the supremum norm in $\mathrm{C}^+[-\tau, T]$. Then for any |
| $(x, y) \in \mathrm{C}^+[-\tau, T] \times \mathrm{C}^+[-\tau, T]$ we define |
| $\|(x,y)\| = \|x\|_T + \|y\|_T,$ |
| $\mathrm{C}_T = \mathrm{C}^+[-\tau, T], \quad \mathrm{C}_T^2 = \mathrm{C}^+[-\tau, T] \times \mathrm{C}^+[-\tau, T].$ |
| |
| Members of $\mathrm{C}_T$ will be denoted by $Z_t$ with the meaning that |
| $Z(t+s) = z_t(s)$ *for* $s \in [-\tau, T]$ *and* $t \in [0, T]$. Now for any fixed |
| $P_t = (P_{1t}, P_{2t}) \in \mathrm{C}_T^2$, (2.8) is a linear system of uncoupled Volterra integral equations in B = $(B_1, B_2)$ and hence (2.8) will |
| have a unique continuous solution under the assumption $A_1 - A_3$. let us denote such a solution by |
| $\mathrm{B}(t) = B^T(P_t), \quad t \in [0, T]$ |
| Or |
| $B_1(t) = B_1^T(P_{1t}, P_{2t}), \quad B_2(t) = B_1^T(P_{1t}, P_{2t}), \ t \in [0, T]$ |
| If we supply this solution $(B_1, B_2)$ in (2.6) we find that (2.) will be satisfied if and only if $(P_{1t}, P_{2t})$ is a fixed point of the |
| operator |

$$\Pi(P_t) = \{\Pi_1(P_{1t}, P_{2t}),\ \Pi_2(P_{1t}, P_{2t})\}$$

Where $\pi_1$ *and* $\pi_2$ respectively denote the right sides of (2.6) for *i =1,2* .

It is not difficult to see from the nature of $M_i$ , $L_i$ and the hypotheses on $f_i, b_i, \varphi_i$ that

$$B: C_T^2 \to C_T^2 \text{ and } \Pi C_T^2 \to C_T^2$$

(which guarantees the nonnegativity of $P_i$ and $B_i$). We will first show that

There exist positive constants $r^*, t^*$ such that $\pi$ is a mapping of a sphere

$$S_{t*}(\Phi, r^*) = \{ (P_{1t}, P_{2t}) \in C_{t*}^2 \mid \|P_t - \Phi_t\|_{t*} \le r^* \}$$

Into itself and also is a contraction where $\Phi_t$ is defined on [-$\tau$, $t^*$]

By the following:

$$\Phi_t = (\Phi_{1t}, \Phi_{2t});$$

(2.12)

$$\Phi_{it} = \begin{cases} \int_0^\infty \varphi_i(a, s)\, da \text{ for } s \in [-\tau, 0]. \\ \int_0^\infty \varphi_i(a, 0)\, da \text{ for } s \in [0, t^*]. \end{cases}$$

We note that similar to (2.11) one gets the estimates

$$|\Pi_i(P)(t) - \Phi_i(0)| \le \{B_i^* \Phi_i(0) \exp[B_i^* t] + \Phi_i(0) d_i^*\} t$$

Which together with (2.7) and (2.12) lead to

(2.13) $$\|\Pi(P_t) - \Phi_t\|_{t*} \le \left( \sum_{j=1}^2 \|\Phi_{jt}\|_{t*} [B_j^* \exp(B_j^* t) + d_i^*] \quad t^* \right)$$

Thus if we choose,

$$r^* = T \sum_{j=1}^2 \|\Phi_{jt}\|_{t*} [B_j^* \exp(B_j^* T) + d_i^*]$$

Then $\Pi$ is a mapping of $S_{t*}(\Phi, r^*)$ into itself for all $t^* \le T$. we then will have to show that for small $t^* > 0$ , $\Pi$ is a contraction.

Let $P_t^{(1)}$ *and* $P_t^{(2)}$ be arbitrarily chosen in $S_{t*}(\Phi, r^*)$ and consider $|\Pi(P_t^{(1)}) - \Pi(P_t^{(2)})|$ for $t \le t^*$.

Corresponding to the chosen $P_t^{(1)}$ , $P_t^{(2)}$ we let $B^{(1)}(a) = B(P_T^{(1)})$ and $B^{(2)}(a) = B(P_T^{(2)})$ . Then from

(2.6),

$$| \text{[illegible]}(P^{(1)}(t)) - \text{[illegible]}(P^{(2)}(t)| \leq \left| \int_0^t B_1^{(1)}(a)\left\{M_1\left(a,t,P_t^{(1)}\right) - M_1\left(a,t,P_t^{(2)}\right)\right\}da \right|$$

$$+ \left| \int_0^t [B_1^{(1)}(a) - B_2^{(2)}(a)]\, M_1(a,t,P_t^{(2)})da \right|$$

$$+ \left| \int_0^t \varphi_1(a,0)\left[L_1\left(a,t,P_t^{(1)}\right) - L_1\left(a,t,P_t^{(2)}\right)\right]da \right|$$

$$= J_1 + J_2 + J_3 \text{ (say)}$$

using the elementary inequality

$$|e^x - e^y| \leq |x-y| \ for \ x,y \leq 0$$

We derive from (2.5) that

$$\left| M_1(a,t,\ p_t^{(1)}) - M_1(a,t,\ p_t^{(2)}) \right|$$

$$\leq \int_o^{t-a} \left| f_1(s,\ p_t^{(1)}(a+s)) - f_1(s,\ p_t^{(2)}(a+s)) \right| ds$$

$$\leq d_1 \| p_t^{(1)} - p_t^{(2)} \|_{t^*} \text{(using } A_1\text{)}$$

*Now,using (2.10),*

$$\left| J_1 \right| \leq \int_o^a \beta_1^* \Phi_1(0)\, e^{\delta 1^a} \left( \int_o^{t-a} d_1 \| P_t^{(1)} - P_t^{(2)} \|_{t^*} ds \right) da$$

$$\leq d_1 \beta_1^* \Phi_1(0)\, t^2 \| P_t^{(1)} - P_t^{(2)} \|_{t^*} e^{\delta 1^t/2}$$

In a similar way

$$\left| J_3 \right| \leq \int_o^\infty \varphi(a,o) \left[ \int_o^t \left| f_1(a+s, P_t^{(1)}(s) - f_1(a+s, P_t^{(2)}(s)) \right| ds \right] da$$

$$\leq d_1 \Phi_1(0) \| P_t^{(1)} - P_t^{(2)} \|_t$$

Now we have to estimate $\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right|$ in order to estimate $J_2$, from (2.8),

$$B_1^{(1)}(a) - B_1^{(2)}(a)$$

$$= \int_o^a \left[B_1^{(1)}(\sigma) - B_1^{(2)}(\sigma) b_1\left(a-\sigma, P_t^{(1)}(a)\right) M_1(\sigma, a, P_t^{(1)}(a) d\sigma + \int_o^a \left[B_1^{(2)}(\sigma)\right.\right.$$

$$\mathrm{X}\left[b_1(a-\sigma, P_t^{(1)}(a) M_1(\sigma, a, P_t^{(1)}(a) - b_1\left(a-\sigma, P_t^{(1)}(a) M_1\left(\sigma, a, P_t^{(2)}(a)\right)\right)\right] d\sigma$$

$$+\int_o^\infty \varphi(\sigma, o)\left[b_1\left(\sigma + a, P_t^{(1)}(a)\right) - b_1\left(\sigma + a P_t^{(2)}(a)\right)\right], L_1\left(\sigma, a, P_t^{(1)}(a)\right) d\sigma$$

$$+\int_o^\infty \varphi(\sigma, o) b_1\left(\sigma + a, P_t^{(2)}(a)\right)\left[L_1\left(\sigma, a, P_t^{(1)}(a)\right) - L_1\left(\sigma, a, P_t^{(2)}(a)\right)\right] d\sigma$$

$$= R_1 + R_2 + R_3 + R_4 \text{ (say)}.$$

By the uniform Lipschitz continuity of $b_1$ in $A_2$ we get

$$\left| b_1\left(a+\sigma, P_t^{(1)}(a)\right) - b_1\left(\sigma + a, P_t^{(2)}(a)\right) \right| \leq B_1 \, /\!/ \, P_t^{(1)} - P_t^{(2)} /\!/ \text{ for } a \, \varepsilon [o, t^*]$$

And hence we have

$$|R_3(a) + R_4(a)| \leq \beta_1 \Phi_1(0) || P_t^{(1)} - P_t^{(2)} ||_{t^*} + \Phi_1(0) \beta_1 d_1 || P_t^{(1)} - P_t^{(2)} ||_{t^*}^{*a},$$

$$|R_2(a)| = [\int_0^a [B_1^{(2)}(\sigma)\{b_1(a-\sigma, P_t^{(1)}(a)) - b_1(a-\sigma, P_t^{(2)}(a))\} M_1(\sigma, a, P_t^{(1)}(a))$$

$$+ b_1(a, -\sigma, P_t^{(2)}(a))\{M_1(\sigma, a, P_t^{(1)}(a)) - M_1(a-\sigma, P_t^{(2)}(a))\}] d\sigma]$$

$$+ b_1\left(a - 0, P_t^{(2)}(a)\right)\{M_1(\sigma, a, P_t^{(1)}(a)) - M_1(a-\sigma, P_t^{(2)}(a))\}] d\sigma]$$

$$\leq \int_0^a \left[\beta_1^* \Phi_1(0) \, \mathrm{e}^{\delta 1^a} \beta_1 \, /\!/ \, P_t^{(1)} - P_t^{(2)} /\!/_{t^*} + \beta_1 \mathrm{d}_1 \, /\!/ \, P_t^{(1)} - P_t^{(2)} /\!/_{t^*} (a-o)\right] do$$

$$\leq \beta_1^* \Phi_1(0) \, \mathrm{e}^{\delta 1^a} /\!/ \, P_t^{(1)} - P_t^{(2)} /\!/_{t^*} \{\beta_1 a + d_1 a^2/2\}$$

Thus

$$\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right|$$

$$\leq \int_o^a \left| B_1^{(1)}(\sigma)\text{-}B_1^{(2)}(\sigma) \right| B_1^{*d}\sigma$$

$$+ \Phi_1(0) /\!/ P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} \left\{ B_1 + B_2\, d_1\, a + B_1^{*} \left( B_2 a + d_1 a^2/2 \right) e^{\delta 1^a} \right\}$$

$$\leq \beta_1^* \int_o^a \left| B_1^{(1)}(\sigma)\text{-} B_1^{(2)}(\sigma) \right| \text{do} + \left( r^* + /\!/ \Phi_t /\!/_{t*} \right) \| P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} g_1^{(a)}$$

Where $g_1(a) \rightarrow \beta_1$ as $a \rightarrow 0$. By Gronwall's inequality it will follow that

$$\left| B_1^{(1)}(a)\text{-}B_1^{(2)}(a) \right| \leq \left( r^* + /\!/ \Phi_t /\!/_{t*} \right) /\!/ P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} \{ g_1(a) + 0(a) \}$$

Using this estimate for $R_1(a)$ we get

$$|R_2(a)| \leq \left| \int_0^a [B_1^{(2)}(\sigma)\, B_1^{(2)}(\sigma),]\, b_1(a)) \text{-} b_1(a\text{-}\sigma, P_t^{(1)}(a))\, M_1(\sigma, a, P_t^{(1)}(a)) do \right|$$

$$\left| B_1^{*} \left( r^* + /\!/ \Phi_t /\!/_{t*} \right) /\!/ P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} 0(a) as\ a \rightarrow 0 \right.$$

*Thus*

$$\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right| \leq (r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ P_t^{(1)} - P_t^{(2)} /\!/_t o(1) as \quad a \rightarrow o$$

And hence

$$\left| \pi_1(P_t^{(1)} - \pi_1\left(P_t^{(2)}\right) \right| \leq ((r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ P_t^{(1)} - P_t^{(2)} /\!/_t o(t) as \quad t \rightarrow o$$

Which together with a similar bound for $\left| \pi_2(P_t^{(1)} - \pi_2\left(P_t^{(2)}\right) \right|$ leads to

$$\| \pi\, (P_t^{(1)} - \pi\, \left(P_t^{(2)}\right) /\!/ \leq (r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ P_t^{(1)} - P_t^{(2)} /\!/_{t*} o(t^*)$$

As $t^* \rightarrow o$ showing that if $t^*$ is small enough, the mapping $\pi$ on

$s_{t^*}(\Phi, r^*)$ is a contraction. Hence the system (2.6)-(2,8) has a unique solution for $o \leq t \leq t^* for\ t^*$ sufficiently small and positive

Now if [0,T] is any finite interval on $R^+$, we can consider [0,T] as a finite union of intervals of length less than or equal to $t^*$ Since the estimate (2.11)is valid for any finite interval we can extend the solution of (2,6)-(2.8) form $[0, t^*]$ to [0,T] and this completes the proof,

3. Stationary age distributions and their linear stability we will now establish the existence of time independent solution

$P_t^*, P_t^*$ Of (1.1)-(1.6): such solutions satisfy

(3.1) $\frac{dp\,P_2^*(a)}{da}+f_i(\mathrm{a}, P_1^*, P_2^*)\,P_i^*(a) = 0,\ \ i = 1,2, a > 0,$

(3.2) $P_i^* = \int_o^\infty P_i^*(a)da, i = 1,\ 2,$

(3,3) $P_i^*=\int_o^\infty b_i(a,\ P_1^*,\ P_2^*)P_i^*(a)da,\ i = 1.\ \ 2$

All solutions of (3.1)are of the from

(3.4) $P_i^*(a) = P_i^*(o)\exp\left[-\int_o^\infty f_i(s, P_1^*, P_2^*)ds\right], i = 1,\quad 2, a \geq 0,$

Which together with (3.3) lead to

$1=\int_o^\infty b_1(a, P_1^*, P_2^*)\exp\left[-\int_o^a f_1(s, P_1^*, P_2^*)ds\right]da,$

(3.5) $1=\int_o^\infty b_2(a, P_1^*, P_2^*)\exp\left[-\int_o^a f_2(s, P_1^*, P_2^*)ds\right]da$

Thus the existence of stationary solutions of (1.1)-(1.6) reduces to the existence of a pair of positive constants $P_1^*, P_2^*$ satisfying (3.5). We can now prove the following,

THEOREM 3. Assume that $f_1$ and $b_i$ satisfy ( $A_1$ ) and ( $A_2$ )and let

(3.6) $F_{i.}\ (P_1, P_2) = \int_o^\infty b_i(a, P_1, P_2)\exp\left[-\int_o^\infty f_i(s, P_1, P_2)ds\right]da,$

$i$ =1, 2, $P_1 \geq 0$, $P_2 \geq 0$,

**And suppose the following hold:**

(3.7) $A_4$. $F_{i.}\ (0,\ 0) > 1,\ i = 1,2$:

$A_5$. *there exists a positive constant c for which*

$F_1$(C,0)< $F_{2.}$(C,0),

(3.8) $F_1$(0.C)< $F_{2.}$(0,C),

Then there exists a unique pair $(P_1^*, P_2^*)$ of real numbers $P_1^* > o$,

$P_2^* > 0$ such that

(3,9) $F_1(P_1^*, P_2^*) = 1 = F_2(P_1^*, P_2^*)$

And the unique nonegative solution of (3.1)-(3.3) is given by

(3.10) $P_2^*(a) = [P_2^* \exp[-f_i(s, P_1^*, P_2^*)ds])/(\int_o^\infty \exp[-\int_o^a f_i(s, P_1^*, P_2^*)ds]da)$,

$i = 1, 2, \quad a \geq 0,$

Proof Consider the elements of the surfaces defined by

Z= $F_1(P_1, P_2), z = F_2(P_1, P_2), P_1 \geq 0, P_2 \geq 0,$

In the ($P_1$,$P_2$,z) space. By hypothesis ($A_1$) we have $\partial F_i / \partial P_j < 0$ *and* , by ($A_2$), $\partial b_i / \partial P_j \leq 0$ for all

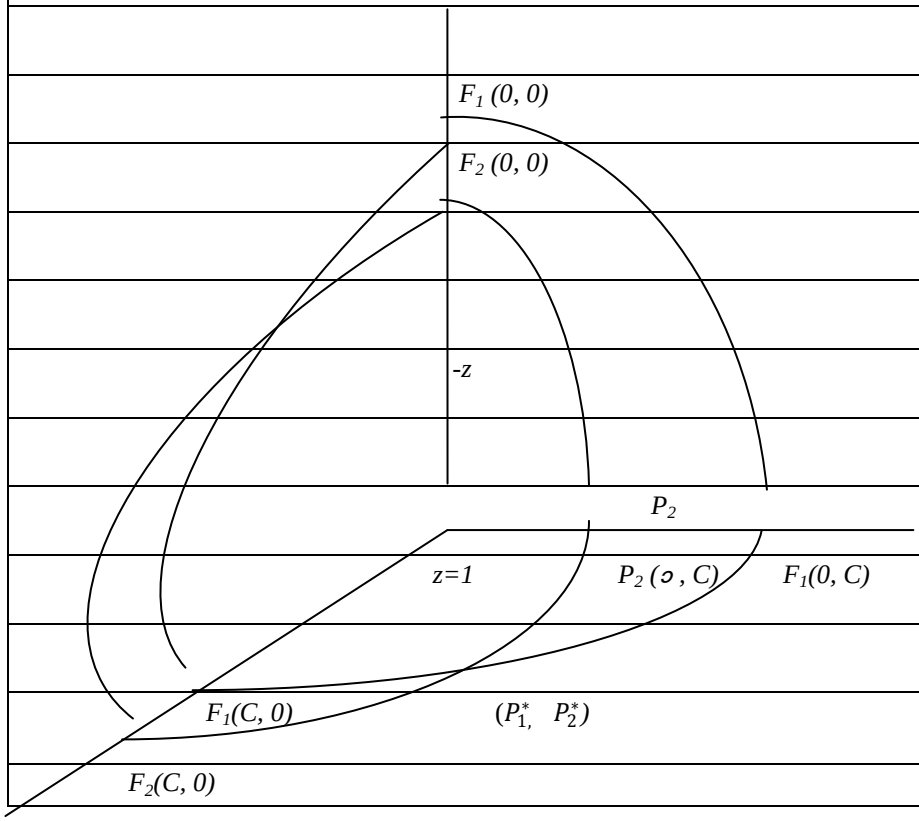
$P_1 \geq 0, P_2 \geq 0, \quad i = 1, 2,$

$j = 1, 2,$ Hence by ($A_4$) the intersections of these surfaces with the plane z=1 define two curves on the z=1 plane connecting the I ines {z=1, $P_1 \geq 0, P_2 \geq 0,$} and {z= $1, P_2 \geq 0, P_1 \geq 0$, These two curves are defined by

$F_1(P_1, P_2 = 1$ *and* $F_2(P_1, P_2) = 1, F_1(P_1, P_2 \geq 0,$

Which by ($A_s$) will intersect at a unique point say $((P_1^*, P_2^*)$ with $(P_1^* > 0, P_2^* > 0$ (see figure). The uniquess of the point $((P_1^*, P_2^*)$ *is a* consequence of the smoothness of the surfaces $z = F_1(P_1, P_2)$ and $z = F_2(P_1, P_2)$

Following is a hypothetical case for Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates attributable or ascribable in the transactions and reactionary potentialities.



(The figure corresponds to the case $F_1(0,0) > F_2(0,0) > 0$. The other possibility $1 < F_1(0,0)$ is treated similarly) It will now

Follow that for such a pair $(p^*_1, p^*_2)$ the unique nonnegative solution of (3.1) - (3.3) is given by (3.10).

Now to examine the linear stability of the stationary age distributions $p^*_1, p^*_2$, we let

(3.11)
$$\begin{cases} p_i(a,t) = p^*_i(a) + u_i(a,t) \\ p_i(t) = p^*_i + p_i(t), \quad i = 1, 2, \quad a > 0, \quad t > 0, \\ p_i(t) = \int_0^\infty u_i(a,t)da, \end{cases}$$

In (1.1)-(1.5) and derive the following variational system after neglecting the nonlinear terms in the perturbations $\boldsymbol{u}_i$ and $\boldsymbol{p}_i$ $(i = 1, 2,)$ (such a procedure of linear stability analysis can be justified As has been done in Gurtin and MacCamy [4].)

(3.12)
$$\begin{cases} Du_i(a,t) = -f_i\left(a, p^*_1, p^*_2\right) u_i(a,t) - p^*_i(a)\left\{\sum_{i=1}^{2} \frac{\partial f_i}{\partial p^*_i} p_i(t)\right\}, \\ u_i(0,t) = \int_0^\infty b_i\left(a, p^*_1, p^*_2\right) u_i(a,t)da + \int_0^\infty p^*_i(a)\left\{\sum_{i=1}^{2} \frac{\partial b_i}{\partial p^*_i} p_i(t)\right\}da \end{cases}$$

Where $i = 1,2$ and (3.13)
$$p_{11}(t) = p_1(t), \quad p_{12}(t) = p_2(t-\tau), \quad p_{21}(t) = p_1(t-\tau), \quad p_{22}(t) = p_2(t)$$

To consider the asymptotic bahaviour as $t \to \infty$ as of solutions of (3.12) we let

(3.14)
$$\begin{cases} u_i(a,t) = u_i(a)\exp[\lambda t], \quad u_i(a) \to 0 \text{ as } a \to \infty, \\ p_i(t) = p^*_i \exp[\lambda t], \\ \qquad\qquad i = 1, 2, \quad a > 0, \quad t > 0, \\ p^*_i = \int_0^\infty u_i(a)da, \end{cases}$$

Using (3.14) in (3.12) - (3.13),

$$\frac{du_i}{da} + \lambda u_i = -f_i\left(a, p^*_1, p^*_2\right) u_i(a) - p^*_i(a)\left\{\sum_{j=1}^{2} \frac{\partial f_i}{\partial p^*_j} p^*_i\right\},$$

(3.15)

$$\mu_i(0) = \int_0^\infty b_i\left(a, p_1^*, p_2^*\right)\mu_i(a)da + \int_0^\infty p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial b_i}{\partial p_j^*}\ p_i^*\right\}da,$$

Where $i = 1,\ 2,$ and

$$p_{11}^* = p_1^*, \quad p_{12}^* = p_2^* \exp\{-\lambda\tau\},$$

$$p_{21}^* = p_1^*, \exp\{-\lambda\tau\}, \ p_{22}^* = p_2^*.$$

If we integrate the first of (3.15) and use the fact that $\mu_i(a) \to 0$ as $a \to \infty$ we can derive that $p_1^*$ and $p_2^*$ are governed by

(3.16) $$\lambda p_1^* - \int_0^\infty p_1^*(a)\left[\left(\frac{\partial b_1}{\partial p_j^*} - \frac{\partial f_1}{\partial p_1^*}\right)p_1^* + \left(\frac{\partial b_1}{\partial p_2^*} - \frac{\partial f_1}{\partial p_2^*}\right)p_2^* e^{-\lambda\tau}\right]da$$

$$= \int_0^\infty \left[b_1\left(a, p_1^*, p_2^*\right) - f_1\left(a, p_1^*, p_2^*\right)\right]\mu_i(a)da,$$

(3.17) $$\lambda p_2^* - \int_0^\infty p_2(a)\left[\left(\frac{\partial b_2}{\partial p_1^*} - \frac{\partial f_2}{\partial p_1^*}\right)p_1^* e^{-\lambda\tau} + \left(\frac{\partial b_2}{\partial p_2^*} - \frac{\partial f_2}{\partial p_2^*}\right)p_2^*\right]da$$

$$= \int_0^\infty \left[b_2\left(a, p_1^*, p_2^*\right) - f_2\left(a, p_1^*, p_2^*\right)\right]\mu_2(a)da,$$

From the first of (3.15) we get

(3.18) $$\mu_i(a) = \left\{\mu_i(0)p_i^*(a)/p_i^*(0)\exp[\lambda a]\right\}$$

$$- p_i^*(a)\exp[-\lambda a]\int_0^a \left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right\}e^{\lambda s}ds, \quad i = 1, 2,$$

Which by the second of (3.15) becomes,

(3.19) $$\mu_i(a)\left[p_i^*(a)/p_i^*(0)c_i\exp(\lambda a)\right]\left[\int_0^a p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right\}da\right.$$

$$\left. - \int_0^\infty b_i\left(a, p_1^*, p_2^*\right)p_i^*(a)\,e^{-\lambda a}\left\{\int_0^a\left(\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right)e^{\lambda s}ds\right\}da\right]$$

$$- p_i^*(a)\,e^{-\lambda a}\int_0^a\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right\}e^{\lambda s}ds, i = 1, 2,$$

where

(3.20) $$c_i = 11 - \frac{1}{p_i^*(0)} \int_0^\infty b_i\left(a, p_1^*, p_2^*\right) p_i^*(a) e^{-\lambda a} da, \quad i = 1, 2,$$

Now supplying $\Phi_1(a)$ and $\Phi_2(a)$ from (3.19) in (3.16)-(3.17) we derive that

$$\left(\lambda - A_{11}\right) p_1^* - A_{12} p_2^* \exp\left(-\lambda\tau\right) = g_{11}(\lambda) p_1^* + g_{12}(\lambda) p_2^* \exp(-\lambda\tau),$$

(3.21)

$$-A_{21} p_1^* \exp\left(-\lambda\tau\right) + \left(\lambda - A_{22}\right) p_2^* = g_{21}(\lambda) p_1^* \exp\left(-\lambda\tau\right) + g_{22}(\lambda) p_2^*$$

Where

(3.22) $$A_{i\Phi} = \int_0^\infty p_i^*(a) \left[\frac{\partial b_i}{\partial p_\Phi^*} - \frac{\partial f_i}{\partial p_\Phi^*}\right] da,$$

(3.23) $$g_{i\Phi}(\lambda) = \left[p_i^*(0) c_i\right]^{-1} \left\{\int_0^\infty \sigma_i(a) p_i^*(a) e^{-\lambda a} da\right\} \Phi \left\{\int_0^\infty \left[p_i^*(a) \frac{\partial b_i}{\partial p_\Phi^*} - b_i\left(a, p_1^*, p_2^*\right) p_i^*(a) e^{-\lambda a} \left(\int_0^a \frac{\partial f_i}{\partial p_\Phi^*} e^{\lambda s} ds\right)\right] da\right\}$$

$$- \int_0^\infty \sigma_i(a) p_i^*(a) e^{-\lambda a} \left(\int_0^a \frac{\partial f_i}{\partial p_\Phi^*} e^{\lambda s} ds\right) da, \quad i, \Phi = 1, 2,$$

(3.24) $$\sigma_i(a) = b_i\left(a, p_1^*, p_2^*\right) - f_i\left(a, p_1^*, p_2^*\right), \quad i, = 1, 2.$$

To solve the linear perturbational system (3.12) it is enough to find the constants $\lambda, p_1^*, p_2^*$ satisfying (3.21) since one can then use (3.19) to find the perturbations from (3.14). A nontrivial solution $\left(p_1^*, p_2^*\right)$ for (3.21) will exist if and only if $\lambda$ is a root of the equation $et \begin{bmatrix} \lambda - \left(A_{11} + g_{11}(\lambda)\right) & -\left(A_{12} + g_{12}\right) e^{-\lambda\tau} \\ -\left(A_{21} + g_{21}\right) e^{-\lambda\tau} & \lambda - \left(A_{22} + g_{22}(\lambda)\right) \end{bmatrix} = 0$

Or equivalently

(3.25) $$\lambda^2 - \lambda\left(A_{11} + A_{22}\right) + A_{11}A_{22} - A_{12}A_{21} \exp\left[-2\lambda\tau\right] = \lambda s_1(\lambda) + s_2(\lambda, \tau)$$

Where $S_1(\lambda) = g_{11}(\lambda) + g_{22}(\lambda)$ ,

$S_2(\lambda, \tau) = \left[A_{21} g_{12}(\lambda) + A_{12} g_{21}(\lambda) + g_{12}(\lambda) g_{21}(\lambda)\right] \exp\left[-2\lambda\tau\right]$

$$-A_{11} g_{22}(\lambda) - A_{22} g_{11}(\lambda) - g_{11}(\lambda) g_{22}(\lambda)$$

The following result establishes the linear local asymptotic Stability of the stationary age distributions $p_1^*(a)$ and $p_2^*(a)$

THEOREM 4 Assume the following:

$A_6.\ 0 < \beta_i^* < d_i^*\ (i,\ =1,\ 2.)\ \ (see(A_1)\ and\ (A_2))$;

$A_7.\ \{A_{11} + A_{22} + \| S_1(0) \|\}^2 < 4\{(A_{11}A_{22} - A_{12}A_{21}) - \| S_2(0) \|\}$,

$$\| S_1(0) \| = |g_{11}(0)| + |g_{22}(0)|,$$

$$\| S_2(0) \| = |A_{11}\, g_{12}(0)| + |A_{12}\, g_{21}(0)| + |g_{12}(0)\, g_{21}(0)|$$

$$+ |A_{11}\, g_{22}(0)| + |A_{22}\, g_{11}(0)| + |g_{11}(0)\, g_{22}(0)|$$

Then all the roots of (3.25) have negative real parts and hence the

perturbations $u_i\ (a,\ t) = \phi(a) \exp[\lambda t] \to 0\ as\ \ t \to \infty$ for $i = 1, 2\ and\ a \geq 0$

Proof . Define $F_1(\lambda, \tau)$ and $F_2(\lambda, \tau)$ as follows:

$$F_1(\lambda, \tau) = \lambda^2 - \lambda(A_{11} + A_{22}) + A_{11}A_{22} - A_{12}A_{21} \exp[-2\lambda\tau];$$

$$F_2(\lambda, \tau) = \lambda S_1(\lambda) + S_2(\lambda, \tau),$$

Using $(A_6)$ one can show that for $\lambda$ with re $(\lambda) \geq 0$, $c_1$ and $c_2$ are positive and bounded away from zero, Now since $A_{ij} > 0$ and $\tau \geq 0$ we find from the nature of the dependence of $g_{ij}$ on $\lambda$ that for all $\lambda\ \ \mu \pm iw$ with $\mu \geq 0$ and $w \geq 0$ we have

(3.26) $$\left| F_1(\lambda, \tau) - P_2(\lambda, \tau) \right|$$

$$\geq \left| F_1(\lambda, 0) \right| - F_2(\lambda, \tau) \Big|$$

$$\geq |\lambda|^2 - |\lambda| \{A_{11} + A_{22} + \| S_1(0) \|\} + \{A_{11}A_{22} - A_{12}A_{21} - \| S_2(0, \tau) \|\}$$

$$\geq \left[ (|\lambda| - P)^2 + (q) \right]^2$$

Where

$$p = \{A_{11} + A_{22} + \| S_1(0) \|\}/2,$$

$$q = \left\{ 4 |(A_{11}A_{22} - A_{12}A_{21}) - \| S_2(0, T) \| | - |A_{11} + A_{22} + \| S_1(0) \| |^2 \right\}/2,$$

It will now follow, from (A) and (3.26), $F_1\ (\lambda, \tau) = F_2(\lambda, \tau)$ connot have roots with zero or positive real parts and this completes the proof.

In conclusion we remark that the condition (3.7) means that the net creation rates of each system is greater than utility; a condition of this type is well known in age wise classification of particles, atoms, molecules, cosmological phenomenon. (3.8) can be interpreted to mean that the intraspecific inhibition is higher than that of the

interspecific interaction. The conditions of Theorem 4 are analytical and do not lend themselves for any worthwhile interpretation in terms of the model parameters. So also an interpretation can be recorded as a remark in respect of destruction rates of particles, properties, elements, or projections.

Note: The model is applicable to each and every module. As stated in the foregoing, there are possibilities of situation where an account itself be a receiver and it is also the one to which debit is raises. Such examples are legion in all the fields of finance, economics, political science and physics and chemistry, and of course mathematics. Classification based on time is the preeminently and predominantly suited premise for such model. We repeat that the model is to be read with each module We will expatiate the importance and implicational ramifications as we proceed with the work. Assortment-blend, collection-combinations, combo-gallimaufry, goulash-hash, hodgepodge-medley, mishmash-mixture, motley-pastiche patchwork-salmagundi, soup-stew, mixed bag-mélange, olio- olio podrida, assortment variety, Array choice, collection combination, combo diversity, hodgepodge jumble, kind medley, miscellany mishmash, mixed bag mixture, mélange potpourri, selection sort, brew concoction, beverage blend, broth compound, distillation drink, fermentation hash, hodgepodge infusion, instillation liquor, mélange miscellany, mishmash mixture, mélange potpourri, borscht bouillon, bowl brew, chowder concoction, consommé decoction, gumbo hodge-podge, olio porridge, potage potpourri, pottage puree, splash stock, vichyssoise hodgepodge, mixture- mess, collection combination, goulash hash, jumble medley, miscellany mishmash, mixed bag mélange, olio patchwork, potpourri salmagundi, jumble clutter, disarrangement disarray, disorder farrago, gallimaufry garbage, goulash hash, litter medley, mess miscellany, mishmash mixture, muddle mélange, olio pastiche, patchwork potpourri, salmagundi scramble, shuffle snarl, tangle tumble, medley miscellany, assortment brew, collection combo, conglomeration farrago, hodgepodge jumble, melee mingling, mishmash mixed bag, mixture mélange, pasticcio pastiche, patchwork potpourri, salmagundi variety, mélange mixture, assortment combo, farrago gallimaufry, pasticcio pastiche, patchwork potpourri, salmagundi soup, stew miscellany, accumulation aggregation, anthology assortment, brew collectanea, combination combo, compilation conglomeration, accumulation diversity, farrago gallimaufry, melee mishmash, mix mixed bag, mixture muddle, mélange odds and ends, olio pasticcio, pastiche patchwork, potpourri salad, salmagundi smorgasbord, stew variety, mishmash jumble, medley mess, miscellany mixed bag, mixture mélange, olio patchwork, potpourri salmagundi, mixed bag array, assortment choice, collection combination, combo conglomerate, diversity hodgepodge, jumble medley, mélange miscellany, mishmash mixture, montage mélange, odds and ends, patchwork potpourri, smorgasbord sundries, mix/mixture assortment, combination, admixture adulteration, alloy amalgam, amalgamation assimilation, association batter, blend brew, combine combo, commixture composite, compound concoction, confection conglomeration, cross crossing dough fusion, goulash grab bag, hodgepodge hybrid, hybridization incorporation, infiltration interfusion, jumble mash, medley merger, mingling miscellany, mishmash mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, odds and ends assortment, bits and pieces, debris etcetera, hodgepodge jumble, knickknacks leavings, leftovers litter, medley melee, miscellaneous paraphernalia, miscellany mishmash, mixed bag motley, mélange notions, novelties oddments, olio particles, patchwork potpourri, rags refuse, remainder remains, remnants rest, rummage scraps, sundries sundry items, compilation mishmash, paste-up patchwork, potpourri reappropriation, reproduction synthesis, conglomeration cross, crossing dough, fusion goulash, interfusion jumble, mash medley, merger mingling, miscellany mishmash, mix mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, colluviums gallimaufry, hash mixture, pasticcio pastiche, conglomeration cross section, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multiformity multiplicity, mélange potpourri, range shift, variegation variousness, divergencies diverseness, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multifariousness multiformity,

mélange potpourri, range shift, soup stew, variance variation, variegation variousness, oddment odds and ends, assortment bits, polymorphism assortment, change collection, combo conglomeration, cross section departure, discrepancy disparateness, divergencies diverseness, diversification diversity, fluctuation heterogeneity, heterogeneousness incongruity, intermixture many-sidedness, medley miscellaneousness, miscellany mishmash, mixed bag mixture, notions novelties, oddment oddments, olio particles, patchwork potpourri of the particles, interactions, accounts, transactions in the cosmos. The wisdom and understanding of our own divine essence is like a beaconing begorra, bureaucratic burdock, befuddling binary and digital digamma, that invites us inward toward ever deeper connection with all around us. Yet rather than look within, most in this world prefer to look outward for divine wisdom, thus sustaining a dependence upon a vast hierarchy that stretches between the individual and the Divine. In all our wanderings away from the Divine, humankind has obscured, obfuscated, obliterated, adjured, abnegated and abjured its most compelling features through a persistent belief systems in limitations arising out of the controlling mechanisms of this hierarchy of sub ordinates, superordinates and coordinates. Has my consciousness changed? Why that doubt? Twist and shake, ponderous novelty themes, pathos, poignancy in writing in different styles, accord-coetaneousness, coevality-coincidence, concurrence-conformity, conjunction-contemporaneousness, harmony, order, peace, simultaneousness- synchronicity exciting optimism of youth, power obmbudsman,pulsations and throbbing of pop, hard luck stories of people, modernity and mischief, grimalkins of rurality, wobbling boats of lovers, saxophone in mouths all point to newer value orientations. The Divine dances and the “Grand design “is perpetrated outside of the confines of any hierarchical structures, be it feudal, organizational. It is complete within itself, and has a singular purpose of demonstrating the collective potential of all life within the universe. It is the archetype of perfection. It is the standard bearer of each soul's innate design and ultimate destiny. The essence of the Divine is far beyond mental conception, yet humanity's tendency is to resort to the limiting language of the hierarchical paradigm acceding- accession, accommodation- accord, accordance, acknowledging, adjustment-affiliation, affinity-alliance, amity-approving, arbitration-arrangement, assenting, authorizing, bargaining-compatibility, compliance--compromise, concert-concession, conformity-congruity, consistency-correspondence, endorsing, granting-harmony in conformity to pen chance ,predilections of the protagonist to define and understand it. Searching for Connection and Wholeness calls for not lumping elements like me made of clay feet and weak knees, owing allegiance to the fringe, without courante and courage of conviction and will for vindication, visual restlessness and exaggerated narratives of self to its own, exciting potentialities available, terrestrial passion, and skew of scams all bear an infallible demonstration for new order of consciousness. Theorists have posited a virtual smorgasbord of objects to account for dark matter, although many of them have fallen prey to observational constraints. As leading possible candidates for baryonic dark matter, there are black holes (large and small), brown dwarfs (stars too cold and faint to radiate), sun-size MACHOs, cold gas, dark galaxies and dark clusters, to name only a few. The range of particles that could constitute nonbaryonic dark matter is limited only slightly by theorists' imaginations. The particles include photinos, neutrinos, gravitinos, axions and magnetic monopoles, among many others. Of these, researchers have detected only neutrinos--and whether neutrinos have any mass remains unknown. Experiments are under way to detect other exotic particles. If they exist, and if one has a mass in the correct range, then that particle might pervade the universe and constitute dark matter. But these are very large "ifs." To a great extent, the details of the evolution of galaxies and clusters depend on properties of dark matter. Without knowing those properties, it is difficult to explain how galaxies evolved into the structures observed today. As knowledge of the early universe deepens, I remain optimistic that we will soon know much more about both galaxy formation and dark matter. What we fail to see with our eyes, or detectors, we can occasionally see with our minds, aided by computer graphics. Computers now play a key role in the search for dark matter. Historically, astronomers have focused on observations; now the field has evolved into an experimental science. Today's astronomical experimenters sit neither at lab benches nor at telescopes but at computer terminals. They scrutinize cosmic simulations in which tens of thousands of points, representing stars, gas and dark matter; interact gravitationally over a galaxy's lifetime. A cosmologist can tweak a simulation by adjusting the parameters of dark matter and then watch what happens as virtual galaxies evolve in isolation or in a more realistic, crowded universe.

Old consciousness will not do. Newer one has to be imbibed and cultivated. When people are unaware of their inner wisdom and wholeness, they tend to search for order and security outside themselves. That is my personal experience. There is always search for sanctity of emotion, aesthetic expression, pursuit of harmony, unraveling realms, syncopated ending and stillness of mind and shiver of heart at the mirror broken somewhere are the strains of newer wisdom if it can be called one. Same questions shall be asked and same shall be answer in the similar vein. Uncertain of their place within the hierarchical order of the world, they define themselves based upon their insecurities. Individuals thus become only pieces of their wholeness and like shards of glass from a beautiful vase; they bear little resemblance to their aggregate beauty. Within the hierarchy, many in high positions of power have taken advantage of our collective insecurities in attempting to guide the development of all humankind. They have obscured the direct connection between the individual and the Divine through a variety of means designed to intercede between our inner essence and our divine source. Each individual can come to know them to be free of all forms of institutional and organisational and societal hierarchical control. This is not to imply that we should not trust others or join together in bonds of friendship and community. It is simply a reminder that relative truth is constantly shifting in the hands of those who desire to control. And even when the motive for their controlling behaviors may be of good will, it is still a form of control. When the revealers of "truth" within the hierarchy withhold and suppress information, they are usually positioning themselves to maintain or acquire more power rather than to disseminate empowerment to all. The desire for connection and wholeness is a fuel that drives us to seek out and explore the hierarchical paradigm. This inner longing provides us with the motivation to seek help and guidance from a specific group within the hierarchy, and in so doing, cultivate a sense of belonging and connection. Furthermore, the hierarchical paradigm is a stage whereby we develop a sense of connection to some grand, encompassing vision. This is why the hierarchy nurtures prophets who point toward a greater vision. Spiritual leaders are able to peer deeply beneath the surface reality of life and experience how intricately connected every life form is, and how the composite of all life is intelligent far beyond measure. These visionary leaders can thus interpret reality through their personal abilities to perceive and express life's dimensional depth and limitless intelligence. Yet no one is able to articulate life's full dimensional depth and breadth with the tools of language. They can only, at best, describe their interpretation or their impressions. In actuality, all of us are able to peer beneath the surface reality of life and perceive a unique vision of the universe. We require only time and intention to develop our own interpretations. And this is precisely what many great spiritual leaders have taught. Life's deeper essence is not an absolute to be experienced by the chosen few, but an evolving, dynamic intelligence that wears as many faces as there are life forms. No group or people have the exclusive portal into the universe by which the Divine expresses itself in all its majesty. The portal is open and available to all, because the Divine is within all things. Those recognized as great prophets each produced a vision of the universe beyond what was currently defined by the hierarchy. Because their interpretations were articulated with authority and depth of insight, they became a target of debate among various groups in the hierarchy. This debate then created a polarity of belief. A sympathetic constituency emerged to defend and embellish their leader's interpretation, while established groups held it in contempt of previously held beliefs. Invariably, the leader's vision became confined and shaped into dogma by followers who desired to create a new religion or sect. Thus, this infusion of fresh insight quietly receded into the hands of the hierarchy, where its deeper meaning was obscured by the very fact that it was incorporated into a rigid structure that both protected and promoted it. A new paradigm is emerging that promotes a clear connection of individual consciousness to the compelling features of the Divine without the intervention of a hierarchy. This is when the fables and myths of history step into the light and become known as they were originally intended. This is the time when language will be transformed into a new form of communication that breaks down all barriers of control. Personal transformation, through the awakening of inner wisdom and understanding, is the pathway into wholeness. The transformation paradigm is initiated simply by the recognition that rather than the dependency-inducing ways of the hierarchical paradigm, there are accelerated, independent pathways that bypass the hierarchy and lead to self-mastery. These new pathways lead to the divine wisdom and understanding that is present within all of us. This wisdom can be accessed through the practice of three principles of transformation: seeing the Divine in

all, nurturance of life, and gratitude. The application of these life principles disengages individuals from the controlling elements of the hierarchy, thus initiating the transformation experience. As there are relative truths, there are relative freedoms. As individuals evolve through the hierarchical process, an ever-increasing sense of freedom is gained, yet external forces continue to exert control through limiting language, confining belief systems, and more. These controlling influences lead to anatomical architecture, body, build, cagey carcass, construction, enclosure, fabric, flounce, form, framework, fringe, groundwork, hem, mount, mounting outline, physique, scaffold, scaffolding scheme, setting shell, stage, structure, support, system, trim, trimming, truss- valance continued reliance upon the hierarchy as it unceasingly attempts to impose a sense of inequality between us. The underlying equation of the hierarchical paradigm is: individual + hierarchy = God connection. In the case of personal transformation it is: individual + inner wisdom and understanding divine equality with all. The time has come to integrate the dominant paradigm of the hierarchy with the liberating transformation paradigm. That shall redeem all promises made to oneself. There shall be a Matrix force or balancing force that shall remain UN satisfies until conservativeness is reached. This integration occurs naturally once we have fully explored the two paradigms and develop a synthesis paradigm whereby transformation is attained by searching for connection and wholeness through our own inner wisdom and understanding. It is this combination of release from dependency on the hierarchy and transformation from within that initiates the synthesis paradigm. Once we take responsibility for our own transformation and integration, it does not mean that the hierarchy is to be shunned or avoided. The hierarchy is quite benign as a manipulative force and merely represents one important stage in the journey toward wholeness. What is being set in motion now is the initial preparation for these paradigm shifts. More specifically, these paradigms will be simultaneously played out over the coming years. As always, it will be the choice of each one of us as to which paradigm we embrace in our journey. All of the highest imaginings of humanity are yet unaware of our deepest foundation. We have sought the upper reaches of the building, yet remain unaware of the foundation's design. It is here, at the very core of existence that the Divine is bursting forth with its creative energy while simultaneously reintegrating with its invitation to wholeness. It is here that equality is realized, not in the lofty places of relative truth lodged in the hierarchy, but rather in the deepest part of the foundational plan of life's original source and ultimate destiny. The origin and destiny of existence is the tone of equality in life. Listen for this tone – this frequency of vibration – and follow it back into the very foundation from whence all things arise and return. Consider these words as symbols only. Feel the truth that stands behind these words, and tap into this empowering energy force that reaches out for you. Know it as a tone or vibration – a resonance that waits for you around every corner in which your life will turn. It is a beacon of the Divine gathering itself into the form of language in order to usher you to a place from which you can experience the formless tone of equality – the bypass of limitation. It is the primal language of our divine source that bestows to you the freedom to generate your own deepest beauty in the expression of your highest truth.(Reference: Wikipedia)

## Gopalasamy Model For Procrastinated Conducive, Felicitous Counter Sallies In Two Specified Systems

Abstract

Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of decay rate of perturbations. Sole intention, main objective and primary concern of

the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are counted in the same vein as that of predator prey. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 9. Preliminary remarks

Almost all the equations of population ecology have been derived starting from the following simple format (Note this can be used for the development of the creed and type of systemic types in case of cosmology and astrophysics and fiancé and econometrics. To give an example of finance, we can think of a account of the company giving rise to various accounts of its employees, and that is the contribution of the single account for the development of the total system types of the given configuration). It may be account, particle, atom, molecule, proposition, problem or any such thing. We shall delineate and disseminate information and particulars as and when the model shall be applied to the problem in question.

$$\frac{dN(t)}{dt} = \text{an individual's contribution}$$

*to population growth* $N(t)$

(1.1)

where $N(t)$ denotes the size of a population of a single species (or biomass) with overlapping generations at time $t$, and $dN(t)/dt$ is the rate of change of the population size at time $t$. Subsequently one makes an assumption regarding the factor in the brackets in (1.1); in particular if one assumes that an individual entity's contribution to the population growth is denoted by a function say, f(t, N), defined suitably for all t>0, N then one obtains from (1.1) the so called Kolmogorov formulation in the form

$$\frac{dN(t)}{dt} = f(t, N(t))\, N(t).$$

(1.2)

Various choices of $f$ together with some ecologically plausible assumptions such as the temporal constancy of the environment and density dependent effects in $f$ lead to various equations. For instance if $f(t, N) \equiv r$ (a positive constant) one obtains the Malthusian formulation in the form

$$\frac{dN(t)}{dt} = rN(t)$$

And if one assumes $f(t, N) \equiv r - (r/K) N$ for some positive constants r and K one gets the familiar logistic equation of population biology

$$\frac{dN(t)}{dt} = rN(t) \left\{ 1 - \frac{N(t)}{K} \right\}$$

In order to device models with time delayed effects, such effects are introduced in the average growth rate $F(t, N)$ of (1.2); for instance if $f(t, N) \equiv r - (r/K) N(t - \tau)$ for some constant $\tau > 0$ we obtain Hutchinson's[19]delay loistic equation

$$\left\{ \frac{dN(t)}{dt} = rN(t) \left\{ 1 - \frac{N(t-\tau)}{K} \right\} \right. .$$

(1.3)

If for instance

$$F(t, N) = r - (r/K) N(t) - \int_{-\infty}^{t} k(t - s) F(N(s))\, ds$$

Where $k$ is a piecewise continuous nonnegative valued function such that $\int_0^\infty k(s)\,ds = 1$ with some appropriate function $F$, we will get Volterra's model of a population which pollutes its environment and pollution has an accumulative and toxic effect; such an equation is of the form

$$\frac{dN(t)}{dt} = N(t) = \left\{ r - (r/K) N(t) - \int_{-\infty}^{t} k(t-s) F(N(s))\,ds \right\} \quad (1.4)$$

Which has been studied in different formats by several authors (Cushing [12], Cohen *et al.* [10], and Gopalasamy [16])

One of the significant advantages of deriving various models from the prototype (1.2) is that if the initial; conditions (initial population sizes) are nonnegative then the nonnegativity of the population size $N(t)$ for $t > 0$ follows from the fact that solutions of (1.2) satisfy

$$N(t) = N(0) \exp\left[\int_0^t f(s, N(s))ds\right]$$

Which implies that $N(t) \geq 0$ if $N(0) \geq 0$.

While the time delayed logistic equation (1.3) explains some of the observed oscillations of certain populations, its derivation has not been entirely satisfactory. For instance it has been pointed out by Ricklefs ([27], page 488) that "time lags might be expected to occur primarily in stabilizing processes involving production rather than death because death is an immediate response to environmental change". Recently Cushing [13] has examined the effects of response and maturation delays in a number of single and multi system models; as it has been noted by Cushing [13], the models involving delayed reproduction or not necessarily derivable from the Kolmogorov type systems of the form (1.2) by suitably defining the average growth rate $(1/N)(dN/dt) = f(t, N(t))$. While Cushing [13] has been mostly concerned with the possibility of delay induced oscillations we will consider a class of delays in reproduction and interspecific interaction which maintain the stability of the systems.

Our starting point to drive models with delayed reproduction is the following balance equation which assumes that there is no immigration or emigration. Account, particle, atom, molecule, proposition, problem et al

$$\frac{dN(t)}{dt} = \text{birth rate} - \text{death rate}. \quad (1.5)$$

Now for instance if we consider a population of adult flies then the birth rate at time $t$ depends on the adult population at time $t^{-}\ \tau$ where $\tau$ is the time required for the larvae to become adult flies. If the birth and death rates are governed by density dependent factors then we have from (1.5) the following time delayed model

$$\frac{dN(t)}{dt} = b\big(N(t^{-}\tau)\big) - m(N(t))$$

(1.6)

Where the functions $b$ and $m$ denote the density dependent birth and death rates of Account, particle, atom, molecule, proposition, problem et al respectively. If the time delay is continuously distributed then one can consider in the place of (1.6) an equation of the form

$$\frac{dN(t)}{dt} = b\int_{-\infty}^{t} k\{(t-s)N(s)ds - m(N(t))$$

Or equivalently

$$\left\{\frac{dN(t)}{dt} = b\int_{0}^{\infty} k(s)\right\}N(t-s)ds - m(N(t))$$

(1.7)

Where the delay kernel $k$ denotes the weight of the delayed of the delayed effects.

It is not always obvious that suitable nonnegative initial conditions for (1.6) and (1.7) will imply the nonnegativity of the solutions of (1.6) and (1.7). The purpose of this article is to show that in two species systems of competitive and prey- predator communities, delayed reproduction and delayed interactions donot destabilise the systems if the self- regulating negative feedback effects are sufficiently strong compared to interspecific interactive effects and if the self regulating effects are realised with no time delays. When we say that time delays do not destabilise, we will mean that delays do not render an asymptotically (locally) stable steady state unstable in such a manner that the loss of stability either leads

To a delay – induced bifurcation to persistent and undamped oscillations (as in the case of Hopf- bifurcation) or the relevant linear variational system has unbounded solutions; that is a mathematical instability is not induced by delays in the sense, that the roots of the characteristic equation are pushed to the imaginary axis or to the right of it on the complex plane. However delays can make an otherwise stable system less stable which means that the rate of convergence to zero of solutions of the variational system can decrease with an increase of delay and the variational system can even oscillate (perhaps violently) before converging to the trivial state. Such a reduction in stability can

happen even when all the roots of the characteristic equation have negative real parts. For our purposes such a reduced stability is still stability. In order to see the effects of delays on such less stable systems we will propose below a method of estimating the rate of decay of local perturbations based on the variational system and indirectly relying on the characteristic equation. For a special class of two species models with a single delay we have obtained (Gopalsamy [16]) estimates of decay rates (or return times) usingthe characteristic equation; such a procedure is almost impossible in the case of many species systems with several delays. The method we propose in Section 6 below for the estimate of decay rates can be adapted to *n*- species systems with $n^2$ discrete delays.

Let $x_1(t)$ and $x_2(t)$ denote the population sizes (or biomasses) of two systems competing for a common pool of resources.. Let $b_i$ and $m_i$ $(i = 1, 2)$ denote the respective density dependent birth (potential) and death rates so that in the absence of any time delays the population sizes are governed by (of Account, particle, atom, molecule, proposition, problem et al). In addition as usual, we assume that the system studied follow Haimovici's equations, for which the dove tailing explanation has been appended

$$\left.\begin{aligned}\frac{dx_1(t)}{dt} &= b_1\left(x_1(t)\right) - m_1(x_1(t)), x_2(t))\\ \frac{dx_2(t)}{dt} &= b_2\left(x_2(t)\right) - m_2(x_1(t)), x_2(t))\end{aligned}\right\} \quad (2.1)$$

In order to make the system (2.1) represent a competition system we will make the following assumptions on the birth and death rates; $_1$, $m_1$ $(i = 1, 2)$ are continuous with continuous partial derivatives for all $x_1 \in 0$ *(I = 1, 2);* also we assume (see for instance Rescigno and Richardson [26]) $\frac{\partial b_i}{\partial x_1} > 0, \frac{\partial m_1}{\partial x_j} > 0$ for $x_i > 0$ $(i, j = 1, 2)$, (2.2)

(xiii) $b_1(0) = 0$, $m_1(0, x_2) \equiv 0, b_2(0) = 0$, $m_2(x_1, 0) \equiv 0$, (2.3)For some $x_{1,}^* > 0, x_{2,}^* > 0$ we have

$$\left.\begin{aligned} b_1(x_1^*, -m_1(x_1^*, 0) &= 0,\\ b_2(x_2^*, -m_1(0, x_2^*) &= 0,\end{aligned}\right\} \quad (2.4)$$

(xiv) There exist positive constants $\delta_1$, $\delta_2$ such that

$b_1(\delta_1) - m_1(\delta_1, x_2) < 0$,

$b_2(\delta_2) - m_2(x_1, \delta_2) < 0$, for $x_1 \geq 0$, $x_2 \geq 0$,

(2.5)

(xv) There exists a point $(\alpha, \beta), \alpha > 0, \beta > 0$ such that

$$b_1(\alpha) - m_1(\alpha, \beta) = 0 = b_2(\beta) - m_2(\alpha, \beta).$$

(2.6)

The conditions on $b_i$ in (2.2) mean that the potential (not average) birth rates are density dependent and any negative density dependent feedbacks to the opened or birth rate of Account, particle, atom, molecule, proposition, problem et al rates are included in the destruction rates; the conditions on $m_i$ indicate intraspecific and interspecific confederation needs and necessities. (2.3) implies that (0, 0) is a (trivial) steady state of the system (2.1) as it is in the case of all ecological models; (2.4) means that in the absence of either of the two species the other one has a nontrivial steady state; (2.5) implies that each of the two populations cannot grow unboundedly since when $x_1(t) > \delta_i, \frac{dx_1}{dt} < 0$. (2.6) means that there exists a nontrivial steady state $(\alpha, \beta)$ in the interior of the positive quadrant of the $(x_1, x_2)$- plane. These assumptions are standard and are minimal in the opinion of the author.

An immediate example of (2.1) –(2.6) is the familiar Volterra –Lotka competition system described by

$$\frac{dx_1(t)}{dt} = r_1 x_1(t) - a_{11} x_1^2(t) - a_{12} x_1(t) x_2(t),$$

$$\frac{dx_2(t)}{dt} = r_2 x_2(t) - a_{21} x_1(t) x_2(t) - a_{22} x_2^2(t),$$

(2.7)

Where $r, a_{ij}, (i, j = 1,2)$ are positive constants satisfying the condition

$$\frac{a_{11}}{a_{21}} > \frac{r_1}{r_2} > \frac{a_{12}}{a_{22}}.$$

One of the first questions to be resolved for (2.1) is the following: if $x_i(0) > 0$, $x_2(0) > 0$ does it follow that the corresponding solutions of (2.1) satisfy the condition $x_1(t) \geq 0$, $x_2(t) \geq 0$ for those $t$ for which a solution of (2.1) exists? In the $(x_1, x_2)$ phase plane of (2.1) we have (0, 0), $(x_1^*, 0)$, $(0, x_2^*)$, $(\alpha, \beta)$ as the steady states of (2.1) and also it will follow from (2.3) that the co ordinate axes

$x_1 = 0$ and $x_2 = 0$ are invariant sets f (2.1) ; such invariance of the coordinate axes will imply that if $x_1(0) > 0$, $x_2(0) > 0$ then the corresponding population trajectory $(x_1(t), x_2(t))$ of (2.1) cannot reach the outside of the nonnegative quadrant $x_1 \geq 0$, $x_2 \geq 0$.

Our plan is to derive first a set of sufficient conditions for the coexistence of the two competing species and then in the next section show that under the same conditions the two species can coexist if there are time delays in the reproduction and interspecific interactions. We recall that the two competing species described by (2.1) are said to coexist if and only if the nontrivial steady state $(\alpha, \beta)$ is (at least locally) asymptotically stable. To examine such a stability of $(\alpha, \beta)$ we derive the corresponding linear variational system for the perturbations $X_1, X_2$ where

$$x_1(t) \equiv \alpha + X_1(t), \qquad x_2(t) \equiv \beta + X_2(t)$$

Such a variational system is found to be

$$\left.\begin{aligned} \frac{dX_1(t)}{dt} &= \beta_{11} X_1(t) + \beta_{12} X_2(t), \\ \frac{dX_2(t)}{dt} &= \beta_{21} X_1(t) + \beta_{22} X_2(t), \end{aligned}\right\} \qquad (2.8)$$

Where

$$\beta_{1J} = \left(\frac{\partial b_1}{\partial x_J} - \frac{\partial m_1}{\partial x_j}\right) \text{ evaluated at } (\alpha, \beta) \qquad (i, j = 1,2).$$

(2.9)

The steady state $(\alpha, \beta)$ of (2.1) is locally asymptotically stable if and only if the trivial solution (0, 0) of (2.8) is asymptotically stable. We can derive the following result from elementary considerations.

In the system (2.1)-(2.6) assume the following:

$$\left.\begin{aligned} \frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1}, \\ \frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1}, \end{aligned}\right\} \qquad \text{at } (\alpha, \beta)$$

(2.10)

Then the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (2.1) – (2.6).

The proof follows easily if we note that the characteristic equation associated with the linear system (2.8) is given by

Det $\begin{pmatrix} \lambda-\beta_{11} & -\beta_{12} \\ -\beta_{21} & \lambda-\beta_{22} \end{pmatrix} = 0$

Or equivalently,

$$\lambda^2 - \lambda\,(\beta_{11} + \beta_{22}) + \beta_{11}\beta_{22} - \beta_{12}\beta_{21} = 0. \qquad (2.11)$$

It follows from (2.2) and (2.10) that,

$$\beta_{11} < 0, \beta_{11}\beta_{22} - \beta_{12}\beta_{21} > 0$$

And hence the roots of (2.11) have negative real parts implying the asymptotic stability of the trivial solution of (2.8)

If apply the condition (2.10) to the Volterra-Lotka model (2.7) we note that when a nontrivial steady state say $(\alpha, \beta)$ exists for (2.7), (2.10) leads to

$$a_{11}\alpha > a_{21}\,\beta \; and \; a_{22}\,\beta > a_{12}\,\alpha,$$

Which together imply $a_{11}\,a_{22} > a_{12}a_{21}$. If $(\alpha, \beta)$ exists it is known that $a_{11}\,a_{22} > a_{12}a_{21} > 0$ is a sufficient condition for the (local) asymptotic stability of $(\alpha, \beta)$.

The condition (2.10) which guarantees the coexistence of competing populations in (2.1)-(2.6) simply means that the intraspecific self – regulating feedback effects $\partial m_i / \partial x_1\ (i = 1, 2)$ bare higher than its own positive feedback $\partial b_1 / \partial x_1\ (i = 1, 2)$ together with its negative (competitive) effects on its competitor $\partial m_I / \partial x_J\ (i,\ j = 1, 2; i \neq j)$.

## 10. Coexistence under delayed production and interactions, transactions

We will now establish that two competing species with a sufficiently strong intraspecific regulation can coexist in the presence of delayed reproduction and delayed interspecific interactions. It is well known that if there are delays in the self – regulating intraspecific effects, stationary coexistence may not always be possible (Gopalsamy [16], Gopalsamy and aggarwala [15], Shibata and Saito [28]). Delayed reproduction is more general and widespread in most biological populations. To build a competition model with such delays we proceed as follows.

Let $\tau_{ij}\ (i, j = 1, 2)$ be a set of nonnegative constants with $\tau = \max \tau_{ij}\ , i, j = 1, 2,$ and suppose that the two ompeting species display delayed reproduction and delayed interspecific interactions with no delays in the self – regulating negative feedback effects. Such a competition system in a constant environment can be modeled by the autonomous delay differential system

$$\frac{dx_1(t)}{dt} = b_1\big(x_1(t-\tau_{11})\big) - m_1\big(x_1(t), x_2(t-\tau_{12})\big),$$

$$\frac{dx_2(t)}{dt} = b_2\big(x_2(t-\tau_{22})\big) - m\big(x_1(t-\tau_{21}), x_2(t)\big), \qquad (3.1)$$

Where both the creation and destruction rates and the morality rates $b_1$ and $m_1\ (i = 1, 2)$ satisfy the same conditions (2.2)-(2.6) as in the previous section. Along with (3.1) we will

Suppose that the initial population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) sizes are provided in the form

$$x_i(s) = \phi_i(s) \ \text{⓪}, \qquad s \in [-\tau, 0], \qquad i = 1,2. \tag{3.2}$$

(Note that the specification of $\phi_i$ need not be on the entire interval$[-\tau, 0]$; but such a hypothesis is convenient.) Since (3.1) is not of the standard Kolmogorov-type, we have to establish that the solutions of (3.1) will remain nonnegative for nonnegative initial conditions.

To prove the nonnegativity of solutions of (3.1)-(3.2), let $\left\{\psi_1^{(n)}(s)\right\}, s \in [-\tau, 0], i = 1,2; (n = 1,2, \ldots)$ be a sequence of strictly positive continuous functions such that

$$\lim_{n\to\infty} \psi_1^{(n)}(s), = \psi_1(s), \qquad s \in [-\tau, 0], \qquad i = 1,2 \tag{3.3}$$

Let $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ be the solution of (3.1) corresponding to the initial condition

$$\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \left\{\psi_1^{(n)}(s), \psi_2^{(n)}(s)\right\}, \quad s \in [-\tau, 0],. \tag{3.4}$$

Consider the solution $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ for $t \in [0, \tau]$ (assuming it exists on the whole of $[0, \tau]$). Suppose $x_1^{(n)}$ does not remain nonnegative for all $[0, \tau^*]$; then there exists a $t^* = [0, \tau^*](\tau^* =$ the positive minimum of $\tau_{ij}$ $(i, j = 1, 2))$ for which

$$x_1^{(n)}(t^*) = 0 \ and \ x_1^{(n)}(t)(t^* - \tau_{11})) > 0.$$

It will follow from (3.1) and the positivity of the initial condition that

$$\frac{dx_1^{(n)}(t^*)}{dt} = b_2\left(x_1^{(n)}(t^* - \tau_{11})\right) > 0 \tag{3.5}$$

Where we have used the properties of $b_i$ and $m_i$ from (2.2)-(2.6); it is found that (3.5) contradicts the definition of $t^*$. Thus $x_1^{(n)}(t) > 0$ on $[0, \tau^*]$; similarly $x_2^{(n)}(t) > 0 \ on \ [0, \tau^*]$. We can repeat this procedure for intervals of the form $[\tau^*, 2\tau^*], [2\tau^*, 3\tau^*] \ etc$. Thus it will follow that so long as $(x_1^{(n)}(t), x_2^{(n)}(t))$ exists, we have

$$x_1^{(n)}(t) \geq 0 \ x_2^{(n)}(t) \geq 0.$$

If we now consider the limits as $n \to \infty$, we get

$$\lim_{n\to\infty}\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \{x_1(t), x_2(t)\} \tag{3.2}$$

With $x_1(t)$ ⓪,, $x_2(t)$ ⓪ where $(x_1(t), x_2(t))$ is the solution of (3.1)-(3.2) and this is a consequence of the continuous dependence of solutions of (3.1) on initial conditions (Hale [17], page 41).

A second question for (3.1) with (2.2)-(2.6) is whether solutions of (3.1) exist for all $t > 0$? suppose now that solutions of (3.1) do not exist for all $t$ ; then there exists a $t_1 > 0$ such that at least for one of either of $x_1$ *or* $x_2$ we have

$$\lim_{t \to t_1 -} x_i\ (t) = \infty. \tag{3.7}$$

To be specific let us suppose that $\lim_{t \to t_1 -} x_i\ (t) = \infty$; let $t_2$ be the first instant of time for which $x_1(t_2) = \delta_1, t_2 < t_1$. It will follow from (3.1) that

$$\frac{dx_1(t_2)}{dt} = b_1\big(x_1\ (t_2 - \tau_{11})\big) - m_1\big(x_1(t_2), x_2\ (t_2 - \tau_{12})\big)$$

$$< b_1\big(x_1\ (t_2)\big) - m_1\big(x_1(t_2), x_2\ (t_2 - \tau_{12})\big)$$

$$< b_1(\delta_1) - m_1\big(\delta_1, x\ (t_2 - \tau_{12})\big) < 0 \tag{3.8}$$

Which contradicts the definition of $t_2$. Thus solutions of (3.1), (2.2) – (2.6) exist for all $t$ .

We are now ready to consider the effects of time delays in reproduction and interspecific interactions on the coexistence of the two competing species. The steady state $(\alpha, \beta)$ of (2.1) is again a steady state of (3.1). To examine the (local) asymptotic stability of $(\alpha, \beta)$ in (3.1) we let

$$x_1(t) = \alpha + Y_1(t), \quad x_2(t) = \beta + Y_2(t)$$

In (3.1) and derive that the linear variational system governing the perturbations $Y_1, Y_2$ is given by

$$\frac{dY_1\ (t)}{dt} = -m_{11}Y_1(t) + b_{11}Y_1\ (t - \tau_{11}) - m_{12}Y_2\ (t - \tau_{12}),$$

$$\frac{dY_2\ (t)}{dt} = -m_{22}Y_2(t) + b_{21}Y_1\ (t - \tau_{21}) - m_{22}Y_2\ (t - \tau_{22}),$$

Where

$$m_{ij} = \frac{\partial m_1}{\partial x_j}, \quad b_{ii} = \frac{\partial b_i}{\partial x_i}\ at\ (\alpha, \beta). \tag{3.10}$$

It is well known (Bellman and Cooke [3], page 336) that the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) if the trivial solution of the variational system (3.9) is asymptotically stable. We formulate our first result in the following form whose proof is given in the appendix.

*THEOREM 3.1. Assume that in the two systemic types competition model system (3.1) with delayed density dependent birth and death rates $b_i$ and $m_i$ $(i = 1, 2)$ satisfy the conditions (2.2)-(2.6). Suppose that the conditions (2.10) hold. Then the nontrivial steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) whatever the delays be $\tau_{ij}$ $(i = 1, 2)$.*

Let $y_1(t) and\ y_2(t)$ respectively denote the population sizes (or biomasses) of a prey species and a predator species which feeds exclusively on the single prey species in a temporally constant environment. In the format of (1.5) the following system will describe the rates of change of the two population sizes:

$$\frac{dy_1(t)}{dt} = f_1(y_1(t)) - g_1(y_1(t), y_2(t)),$$
$$\frac{dy_2(t)}{dt} = f_2(y_1(t)) - y_2(t) - g_2(y_2(t)), \quad (4.1)$$

Where $f_i\ and\ g_i\ i = 1,2)$ respectively denote the birth and death rates of the two species; in order to make (4.1) denote the dynamics of a prey – predator species we will assume that $f_i\ and\ g_i$ satisfy the following conditions (see for instance Rescigno and Richardson [26]); $f_i\ and\ g_i$ are continuous with continuous partial derivatives for all $y_i \geq 0\ (i = 1, 2); also\ we\ require$

$$\frac{\partial f_1}{\partial y_1} > 0; \quad \frac{\partial g_1}{\partial y_1} \geq 0; \quad \frac{\partial g_1}{\partial y_2} \geq 0,$$
$$\frac{\partial f_2}{\partial y_1} \leq 0; \quad \frac{\partial g_2}{\partial y_2} \geq 0; \quad \frac{\partial g_2}{\partial y_2} \geq 0, \quad \text{Where } y_1 \geq 0, y_1 \geq 0, \quad (4.2)$$

(xvii)

$$f_1(0) = 0, \quad g_1(0, y_2 \equiv 0,$$
$$f_2(0, y_2) = 0 = f_2(y_1, 0); \quad g_2(0) = 0, \quad (4.3)$$

(xviii) For some constant $y_1^* > 0$ we have

$$f_1(y_1^*) - g_1(y_1^* 0) = 0, \quad (4.4)$$

(xix) There exists a point $(\alpha^*\beta^*), \alpha^* > 0, \beta^* > 0$ such that

$$f_1(\alpha^*) - g_1(\alpha^*, \beta^*) = 0,$$
$$f_2(\alpha^*, \beta^*) - g_2(\beta^*) = 0, \quad (4.5)$$

(xx) There exist positive constants $n_1, n_2$ such that

$$f_1(\eta_1) - g_1(\eta_1, y_2) < 0, y_2 \geq 0,$$
$$f_2(\eta_1, \eta_2) - g_2(\eta_2) < 0 \quad (4.6)$$

The conditions (4.2) mean that the system (4.1) is of the oppugnacy systems type; (4.3) will imply that (0,0) is a trivial steady state of (4.1); (4.4) means that in the absence of the predator, the prey species has a nontrivial (positive) steady state;

(4.5) implies that (4.1) has a nontrivial steady state in the interior of the positive quadrant of the state species of (4.1); (4.6) can be shown (see below) to imply that no credit type and debit type system types can grow unboundedly.

An example of (4.1)-(4.6) is the familiar Volterra-Lotka prey-predator system

$$\frac{dy_l(t)}{dt} = r_l y_l(t) - [a_{ll} y_l^2(t) + a_{12} y_l(t) y_2(t)],$$
$$\frac{dy_2(t)}{dt} = a_{12} y_1(t) y_2(t) - [r_2 y_2(t) + a_{22} y_2^2(t)],$$

Where $r_{l,}\ a_{ij}\ (i,j=1,2)$are positive costants such that

$$a_{21}r_1 - r_2a_{11} > 0.$$

As in the previous case of competition we have to verify that when $y_1(0) > 0, y_2(0) > 0$ the corresponding solution of (4.1) will be nonnegative for $t$ 0 for which the solution exist. Since the co-ordinate axes $y_1 = 0, y_2 = 0$ are invariant sets for (401) by virtue of (4.3), it will follow that the solution of (4.1) starting from the interior of the positive quadrant cannot reach the outside of that quadrant. This is a consequence of the invariance of the co-ordinate axes for (4.1)

To prove that solutions of (4.1)-(4.6) exist for all t 0 let us suppose that $y_l 0\ and\ y_2(0) satisfy\ 0 < y_l\ (0)$ satisfy $0< y_l(0) < \eta_l, 0 < y_2(0) < \eta_2$ ( on the other hand if either one of $y_l(0)$ $\eta_l and\ y_2(0)$ $\eta_l$ or both hold, it will follow that the corresponding $y_l(t)$ will decrese by virtue of (4.6) at least for small $t$ 0, ). Suppose $y_l(t)$ does not exist for all $t$ 0; then for some $t_l < \infty\ lim_{t\to t_l-} y_l(t) = \infty$and let $t_2 > 0$ be the first instant for which $y_1(t_2) = \eta_1$; then we have from (4.1) and (4.6) that

$$\frac{dy_1(t_2)}{dt} = f_1(\eta_1) - g_1(\eta_{1,}\ y_2) < 0\ (assumimg\ y_1(0) < \eta_1)$$

Which contradicts the definition of $t_2$ showing that $y_1(t)$ exists for all $t$ 0; if $y_2(t)$ is not defined for all $t$ 0, then for some $t_3$ $\infty$ we will have $lim_{t-t_3-} y_2(t) = \infty$;now let $t_4$ be the first instant for which $y_2(t_4) = \eta_2$; then from the properties of $f_2$ we have

$$\frac{dy_2(t_4)}{dt} = f_2(y_1(t_4), \eta_2) - g_2(\eta_2)\ f_2(\eta_1, \eta_2) - g_2(\eta_2)$$

$$< 0 (by\ (4.6) assuming\ 0 < y_1(0) < \eta_1, 0 < y_2(0) < \eta_2)$$

Which again contradicts the definition of $t_4$. Thus $y_2(t)$ is defined for all $t$ 0.

We will now derive a set of sufficient condition for the (local) asymptotic stability of the steady state $(\alpha^*, \beta^*)$ of (4.1). in the next section we will show that under the same conditions and in the presence of various delays , the steady state $(\alpha^*, \beta^*)$remains (locally) asymptotically stable.

To examine the asymptotic stability of the nontrivial steady state $(\alpha^*, \beta^*)$ in (4.1) we let

$$y_1(t) = \alpha^* + U_1(t), \qquad y_2(t) = \beta^* + U_2(t)$$

In (4.1) and derive the liner variational system in the perturbations $U_1, U_2$ as follows:

$$\left.\begin{aligned} \frac{dU_1(t)}{dt} &= \frac{\partial f_1}{\partial y_1}U_1(t) - \frac{\partial g_1}{\partial y_1}U_1(t) - \frac{\partial g_1}{\partial y_2}U_2(t), \\ \frac{dU_2(t)}{dt} &= \frac{\partial f_2}{\partial y_1}U_1(t) - \frac{\partial f_2}{\partial y_2}U_2(t) - \frac{\partial g_2}{\partial y_2}U_2(t), \end{aligned}\right\} \quad (4.7)$$

Where all the partial derivatives in (4.7) are evaluated at $(\alpha^*, \beta^*)$. The steady state $(\alpha^*, \beta^*)$ of (4.1) is (locally) asymptotically stable if and only if the trivial solution of (4.7) is asymptotically stable which will be the case if and only if the roots of the characteristic equation

$$\text{Det} \begin{pmatrix} \lambda - (f_{11} - g_{11}) & g_{12} \\ -f_{21} & \lambda - (f_{22} - g_{22}) \end{pmatrix} = 0,$$

Where $f_{ij} = \frac{\partial f_1}{\partial y_l}$, $g_{ij} = \frac{\partial g_1}{\partial y_j}$, at $(\alpha^*, \beta^*)$ (4.8)

$(i,j = 1, 2)$

Have negative real parts. If $\lambda_1, \lambda_2$ are the riits of (4.8) then

$$\lambda_1 + \lambda_2 = (f_{11} - g_{11}) + (f_{22} - g_{22}), \quad (4.9)$$

$$\lambda_1 \lambda_2 = (f_{11} - g_{11})(f_{22} - g_{22}) + f_{21} g_{12}.$$

The proof of the following result is an immediate consequence of (4.7)-(4.9)

THEOREM 4.1 *In the CREDIT-DEBIT system (4.1)-(4.6) assumes that the partial derivatives satisfy the following additional conditions:*

$$\frac{\partial g_1}{\partial y_1} > \frac{\partial f_1}{\partial y_1} + \frac{\partial f_2}{\partial y_1},$$

$$\frac{\partial g_2}{\partial y_2} > \frac{\partial f_2}{\partial y_2} + \frac{\partial g_1}{\partial y_2}, \quad \text{Evaluated at } (\alpha^*, \beta^*), \quad (4.10)$$

*Then the steady state $(\alpha^*, \beta^*)$, is (locally) asymptotically stable for the prey- predator system.*

## 9. Delays in reproduction, predation and gestation

A number of authors have proposed and analysed several models of prey- predator systems with time delays. Most of the models have been derived from the so called models of the Kolmogorov-type f (1.2) with time delays incorporated in the average growth rates of the prey and predator species. Volterra [32] proposed the following model for a prey-predator system. The equations are used for the formulation of Credit-debit system of atoms, accounts, economic variable, cosmological variable, astrophysical; variables, and qubit systems for which the model is being applied.

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \delta_1 N_1(t) - \int_{\infty}^{t} F_1(t-s) N_2(s) ds],$$

(5.1)

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 + \delta_2 N_1(t) - \int_{\infty}^{t} F_2(t-s)N_1(s)ds],$$

Where $r_1, r_2, \delta_1, \delta_2$ *are positive constant and* $F_1, F_2$ is nonnegative continuous delay kernel defined and integrable on $[0, \infty]$. (5.1) has no predator self-regulating negative feedback. Brelot [6] considered a modified format of (5.1) in the form

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \lambda_1 N_1(t) - \delta_1 N_2(t) - \int_{-\infty}^{t} F_1(t-s)N_2(s)ds],$$

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 - \lambda_2 N_2(t) - \delta_2 N_1(t) - \int_{-\infty}^{t} F_2(t-s)N_1(s)ds], \quad (5.2)$$

Which has term representing intraspecific competition among both predator and prey species. Under suitable condition on the various parameters in (5.2) one can show that (5.2) has a (locally) asymptotically stable steady state. A number of intergrodifferential equation models of the types (5.1) and (5.2) have been investigated by Cushing [12]

Starting from Hutchinson's (1948)[19] delay logistic equation, May[23] has proposed the following system

$$\frac{dN_1(t)}{dt} = rN_1(t)\left[1 - \frac{N_1(t-\tau)}{K}\right] - \alpha N_1(t)N_2(t),$$

$$\frac{dN_2(t)}{dt} = -bN_2(t) + \beta N_2(t)N_1(t), \quad (5.3)$$

Where r, $\tau$, K, $\alpha, \beta, b$ are positive constants; (5.3) contains a single discrete delay; one can modify (5.3) and incorporate a continuously distributed delay in it so that (5.3) can be written as (May [23])

$$\frac{dN_1(t)}{dt} = rN_1(t)[1 - \int_{-\infty}^{t} F(t-s)N_1(s)ds] - \alpha N_1(t)N_2(t),$$

$$\frac{dN_2(t)}{dt} = -bN_2(t) + \beta N_1(t)N_2(t). \quad (5.4)$$

The first model of a prey-predator type which departs from the Kolmogorov-type formulation is due to Wangersky and Cunningham [33] who considered the model system

$$\frac{dN_1(t)}{dt} = a_1N_1(t) - b_1N_1^2(t) - c_1N_1(t) - N_2(t),$$

$$\frac{dN_2(t)}{dt} = a_2N_2(t) + c_2N_1(t-\tau)N_2(t-\tau), \quad (5.5)$$

Where $a_1, a_2, b_1, c_1, c_2, \tau$ are positive constants; (5.5) means that a duration of $\tau$ time units elapses when an individual prey is killed and the moment when the corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) Recently Cushing [13] has considered a model of the form

$$\frac{dN_1(t)}{dt} = rN_1(t)\left[1 - \frac{N_1(t)}{k} - \alpha N_2(t)\right],$$

$$\frac{dN_2(t)}{dt} = -d_2^*(t) + bN_1(t)\int_0^\infty \delta(a)N_2(t-a)e^{-d*a}\,da,$$

(5.6)

Where $r, k, \alpha, d^*, b$ are positive constants and $\delta$ is related to the age dependent fecundity of the predator species.

With the models (5.1)-(5.6) in the background let us now consider a time delayed model similar to (4.1) in the following form:

$$\frac{dN_1(t)}{dt} = f_1\big(N_1(t-\tau_{11})\big) - g_1\big(N_1(t), N_2(t-\tau_{12})\big),$$

$$\frac{dN_2(t)}{dt} = f_2\big(N_1(t-\tau_{21})\big), N_2(t-\tau_{12}) - g_1(N_2(t)),$$

(5.7)

Where $\tau_{ij}$ $(i,j=1,2)$ are nonnegative constants with $\tau = \max \tau_{ij}, i,j = 1,2$; the birth and death rate functions $f_1, f_2, g_1, g_2$ satisfy the same conditions as in (4.1)-(4.6). Along with (5.7) we specify the initial conditions as follows:

$$N_1(s) = \phi_1(s) > 0,$$

$$s \in [-\tau, 0].$$

$$N_1(s) = \phi_1(s) > 0,$$

We3 will assume that in addition to (4.2) we have for (5.7) that

$$\frac{\partial f_2}{\partial N_1} > 0, \frac{\partial f_2}{\partial N_2} > 0, f_2(N_2, N_2) > 0 \; whenever \; N_1 > 0, \; N_2 > 0. \; (5.9)$$

It is not difficult to show by the methods previously used for the competition system that whenever $N_1(s) > 0$ and $N_2(s) > 0$ for $s \in [-\tau, 0]$, solution of (5.7) exist for all $t \geq 0$ and remain nonnegative for all $t \geq 0$. also $(\alpha^*, \beta^*)$ of (4.5) is a steady state of (5.7).

We will discuss the form and meaning of the delays in (5.7) in the next section. The effects of the various delays in (5.7) on the (local) asymptotic stability of $(\alpha^*, \beta^*)$ can be derived from the following theorem whose proof is quite similar to that of Theorem 3.1 and hence the details of the proof are omitted.

THEOREM 5.1 *In the time delayed prey-predator system(* corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) *) (5.7) let the conditions (4.2)-(4.6), (5.8) and (5.9) hold for the birth and death rates; furthermore suppose the condition (4.10) hold. Then for all nonnegative delays in (5.7), its nontrivial steady state $(\alpha^*, \beta^*)$ is (locally) asymptotically stable.*

## 6. Estimation of decay rates

We will first derive a convention sufficient condition for the asymptotic stability of the trivial solution of a linear system of the form

$$\left.\begin{aligned} \frac{dY_1(t)}{dt} &= -m_{11} Y_1(t) + b_{11} Y_1(t-\tau_{11}) - m_{12} Y_2(t-\tau_{12}), \\ \frac{dY_2(t)}{dt} &= -m_{22} Y_1(t) + m_{21} Y_1(t-\tau_{21}) - b_{22} Y_2(t-\tau_{22}), \end{aligned}\right\} \; t > 0, \quad (6.1)$$

Where $m_{1J}, b_{i1}$ $(i, j = 1,2)$ are positive constants while $\tau_{1J}(i, j = 1,2)$ are non negative with $\tau = \max(\tau_{1J})(i, j = 1,2$.as is shown in the appendix, the following system dominates (6.1*)*

$$\left.\begin{aligned} \frac{dZ_1(t)}{dt} &= -m_{11} Z_1(t) + b_{11} Z_1(t-\tau_{11}) - m_{12} Z_2(t-\tau_{12}), \\ \frac{dZ_2(t)}{dt} &= -m_{22} Z_1(t) + m_{21} Z_1(t-\tau_{21}) - b_{22} Z_2(t-\tau_{22}), \end{aligned}\right\} \; t > 0, \quad (6.2)$$

$Z_i(s) = |Y_i(s)|, \qquad i = 1,2 \;\; for \; s \leq 0 \;\; for \; which \; Y_i(s) \; is \; defined.$

Thus the asymptotic stability of the trivial solution of (6.2) implies that of (6.1). it is also shown in the appendix that the solutions of (6.2) are nonnegative for all t $\geq 0$. we will use the following vector-matrix notation:

$$Z = (Z_1, Z_2)^{\tau}, \qquad ||Z|| = |Z_1(t)| + |Z_2(t)| = Z_1(t) + Z_2(t)$$

$$B = \begin{bmatrix} b_{11} & m_{12} \\ m_{21} & b_{12} \end{bmatrix}, \quad ||B|| = max \begin{matrix} b_{11} + & m_{21} \\ b_{21} + & m_{12} \end{matrix}, \tag{6.3}$$

$$A = \begin{bmatrix} -m_{11} & 0 \\ 0 & -m_{22} \end{bmatrix}.$$

The logarithmic norm or the measure of any square matrix B, denoted by $\mu(B)$is defined as follows:

$$\mu(B) = \lim_{\theta \to 0+} \frac{/\!/I + \theta B/\!/ - 1}{\theta}$$

($I$ being an identity matrix of the same size as $B$).The matrix measure compatible with the norms in (6.3) is given by

$\mu(B) = \max_J [b_{jj} + \sum_{\substack{i=1 \\ i \neq J}}^{2} | b_{jJ} | ]$ $B=(b_{1J})$.

We refer to Coppel [11]or Vidyasager [31]for a number of properties of matrix measure We have shown in the appendix that the derivative of $H(Z_1,Z_2)(t)$(seeA.10)satisfies the relation

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq [-m_{11}+b_{11}+m_{21}]Z_1(t)+[-m_{22}+b_{22}+m_{12}]Z_2(t)$, and hence we have

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq \{\mu(A) + /\!/B/\!/\}\{ Z_1(t)+ Z_2(t)\}$

If we assume for the matrix norm $\|B\|$ as in (6.3)that

$$\mu(A) + /\!/B/\!/ < 0 \; or \; -\mu(A) > /\!/B/\!/$$

Then from the results of the appendix (Lemmas 5 and 6) it will follow that $Z_1(t)+Z_2(t) \to 0$ as $t \to \infty$ and thus we have the following:

Lemma 6.1 Assume that the Coefficients of the linear system (6.1)satisfy the condition (6.7). Then the trivial solution of(6.1)is asymptotically stable.

When (6.7)holds an estimate for the decay rate of solutions of (6.1)can be obtained as follows A necessary and sufficient condition for the asymptotic stability of the trivial solution is that all the roots of the characteristic equation

Det $G(\lambda,T_{iJ})=0$

Have negative real parts where

$G(\lambda,T_{1J})=[\begin{matrix} \frac{\lambda+m_{11}}{0} & \frac{0}{\lambda+22} \end{matrix}] - b_{11}^{e-\lambda T_{11}}$ $\begin{matrix} -m_{12}^{e-\lambda T_{12}} \\ -m_{21}^{e-\lambda T_{21}} \quad b_{22}^{e-\lambda T_{22}} \end{matrix}$

$=\lambda I - A - B(\lambda, T_{ij})$(say),

A positive number $\mathfrak{a}$(if it exists)is said to be a decay rate of solution of (6.1) if and only if every root of (6.8)satisfies the condition

$\mathrm{Re}(\lambda)+\mathfrak{a} < 0.$

When (6.11)holds the solution of (6.1)will be of the order of $e^{-at}$ and hence the supremum of all such $\mathfrak{a} > 0$ $a$ satisfying (6.11)will be a desirable parameter to estimate However such a supremum of all such (6.11)may not be possible to determine since it has to be done directly from (6.8)-(6,10)(see also Section 3).We will obtain an estimate of $\mathfrak{a}$ (not necessarily the supremum)I (6.11)by the following procedure.

Replace $\lambda$ byz-$\mathfrak{a}$ in (6.8)and define

$G(z-\mathfrak{a}, T_{IJ}) = M(z) = (z-a)I - A - B(z-\mathfrak{a}, T_{IJ}).$

By Lemma 6.1 a Sufficient condition for all roots of M(z)=0 to have negative real parts equivalently to satisfy

$\mathrm{Re}(\lambda) < -a$

is the following :

$-\mu(aI + A) > \|B_1(a, T_{ij})\|,$

Where

$$B_1(a, T_{ij}) = \begin{pmatrix} b_{11}^{\,e-\lambda T_{11}} & -m_{12}^{\,e-\lambda T_{12}} \\ -m_{21}^{\,e-\lambda T_{21}} & b_{22}^{\,e-\lambda T_{22}} \end{pmatrix}$$

Since $\mu(\mathfrak{a} I + A) = \mathfrak{a} + \mu(A)$ we have from (6.13) that

$\mathfrak{a} + \mu(A) + \|B_I(\mathfrak{a}, T_{IJ})\| < 0$

Is a sufficient condition for Re $(\lambda) < -a$If we denote the left side of (6.15)by $\mathfrak{a} f(\mathfrak{a}, T_{Ij})$ we note that

$f(0, T_{Ij}) = \|B\| + \mu(A) < 0$ $if$ (6.7)$holds$;

Al so $\partial f / \partial \mathfrak{a} > 0, f$ is continuous and monotonic in $(\mathfrak{a})$;$\to \infty as\ \mathfrak{a} \to \infty$; thus by the continuity of $f$ on $\mathfrak{a} \geq 0$, there exists an interval of the from $[0, ^{*})$such that

$f(0, T_{Ij}) < 0$ $for\ \mathfrak{a} \in [0, \mathfrak{a}^{*})$

And hence

$f(\mathfrak{a}, T_{Ij}) = 0$

Thus an upper bound $\mathfrak{a}^*$ of $\mathfrak{a}$ for which we have $f(\mathfrak{a}, T_{Ij}) < 0$ $\mathfrak{a} \in [0, \mathfrak{a}^*)$ can be obtained as the unique (by the monotonicity of $f$ on $\mathfrak{a}$) positive root ofmax $b_{11}^{e-\lambda T}{}_{11}$ - $m_{12}^{e-\lambda T}{}_{21}$ + $\mathfrak{a}$-min $m_{11}$ =0

$b_{22}^{e-\lambda T}{}_{22}$ - $m_{21}^{e-\lambda T}{}_{12}$ $m_{22}$

Where the maximum in (6.18) is taken over $b_{11}, b_{22}, m_{12}, m_{21}$ and $T_{ij}$ $(i, = j\ 1,2)$. A numerical procedure can be devised to solve equations of the form (6.18) in concrete cases. If one is interested in a smaller bound on the decay rate it can be obtained as the unique root of

$\| B \| e^{ar} + a + \mu(A) = 0, \tau = \max_{i,J} T_{ij}$. (6.19)

It will follow from (6.19) than an increase in T

Will reduce the value of the root of (6.19). We remark that since our estimate $a^*$ of dacay rate is based on a sufficient condition of the form (6.7), we can conclude that the solotions of (6.1) will converge at least as fast as $e^{-a*t}$ where $a^*$ is the unique root of either of (6.18) or (6.19).

## Implications

Authors have formulated a class of two systems model systems staring from the elementary principal that the rate of change of a population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) equals the creation rate minus the destruction rate. This is a departure from the conventional model building in population ecology where one usually starts from the Malthusian formulation, now known as the Kolmogorov formulation in the nonlinear context. In our analysis of the models proposed we have obtained sufficient conditions for the global existence of nonnegative solutions of the model systems (this is trivial in the Kolmogorov formulation) and the other desirable properties such as the existence of suitable steady states. The two classes of models proposed for competition and debt contain intraspecific self – regulation biologists that each systems is under the action of a self-regulating negative feedback force, opinions and arguments differ regarding the mechanism by which such an effect is realized; we refer for more details on this aspect to a collection of articles compiled by Tamarin [30].

Gopalasamy's main result is the following : if a self –regulating negative feedback mechanism (intraspecific oppugn ant) can act without time delay and if such an effect can keep an ecosystem ( a community of one or more species of populations) asymptotically stable ( at least locally about a non trivial steady state, then time delays of any magnitude in any other interaction involving positive or negative feedbacks cannot destabilize such a stable ecosystem ( in a mathematical sense); that is such delays cannot cause a hopf-type bifurcation leading to persistent and undamped fluctuations. The delay-induced bifurcations to oscillations in models have been investigated by many authors. (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Qualitatively result is not new to the literature the effects of delays on otherwise stable and unstable systems have been considered by Cushing[12]where (pages 35-39)he has shown that an otherwise unstable system can be made

less unstable by delays, Such an effect of delays on unstable systems appear to have been first proposed by Beddinton and May [2]who have also shows that delays can make stable systems less stable leading to a conclusion "time delays are not necessarily destabilizing" It is also shown by Cushing([12], pages 75-103)that while time delays cannot stabilize(in the mathematical sense)an otherwise unstable two species system in competition or mutualism time delays in there species systems can stabilize otherwise unstable systems (even in a mathematical sense)Cushing and Saleem [14]have formulated a model where delays in the positive feedback can have a stabilizing effect: Hale[18]discusses briefly the stabilizing nature of delays and given an example to show that an increase in delays can make an unstable system stable Thus for a better understanding of the effects of delays on otherwise stable systems one has to estimate the rate of decay of perturbations (or the characteristic return time )in terms of the parameters of the system especially when the delays do not cause a mathematical instability This aspect has been investigated by a procedure slightly different from that of the characteristic equation we have proposed (section6) method for the estimation of decay rate of perturbations when there are multiple delays Since many species systems with multiple delays cannot be analysed (due to complexity)by the characteristic equation our methods using comparison lead to easily verifiable sufficient conditions The sufficient conditions obtained are in the form of a diagonal dominance type; these conditions can be weakened to quasi-diagonal dominance type as in Ladde [20], Siljak[20]Anderson [1], Lewis and Anderson [21,22](Ladde's result [20] is delay dependent due to the fact that delay appears in self –regulation also).Verification of quasi –diagonal dominance conditions in arbitrary matrices (except for Metzler matrices )is an impossible task and quasi –diagonal dominance is hard to interpret.

If follows from our discussion that those confederation systems which display delay induced instability leading to permanent and undamped oscillations should have one or more of the following aspects:

(1)time delays are in intraspecific interaction terms (and possibly in other terms also): (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator –population.)

(2) Intraspecific self-regulation terms are absent in density dependent from in one or more populations of the system;

(3) Intraspecific negative feedback effect are not strong enough to maintain the stability of the relevant steady state of the nondelayed system.

We can now infer how the delay induced stability leads to fluctuations; a primary mechanism of such fluctuations can be described at least theoretically as follows : in its attraction of a system to a steady state, intraspecific self-regulation fails to be effective and the system overhoots the steady state and at such higher densities intraspecific competition becomes operationally effective leading the system to understood the steady state and this process is repeated again and again resulting in undamped oscillations about the steady state. It follows from our results that if intraspecific competition is sufficiently strong, the system cannot become oscillatory even when there are time delays, so long as they do not interfere with the strength of the intraspecific self – regulation.

We are thus led to the question of existence and importance of intraspecific competition in population systems. In the introduction to an anthology of articles by several authors, Tamarin [ 30] begins with the statement "all animal populations have the reproductive capacity to increase without limit, yet none di" and in his concluding remarks writes " this leaves us with only one great concept in the area of confederation do main regulation: that all credit-debit are regulated" More specifically this is applicable to all the confederation of (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions,

financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

In his discussion on the effect of competition on the densities of populations, Nicholson ([ 25], $page$ [ 147] ) has reached the conclution "Intraspecific competition automatically regulates the severity of its action to the requirements of each denizen , provided the inherent resistance of the environment is sufficiently low to permit the species to exist." In order to look for a mechanism by which intraspecific competition plays its role of decreasing growth rate as confederation density increases, Calhoun [ 7] set up an experiment with Norway rates (Rattus norvegicus) and found that increased density leads to increased hostility which disturbs the psychology and physiology causing decreased reproductive output. Many other mechanisms by which intraspecific compition self-regulate populations have been put forward by a number of authors (see Tamarin [30]). This is true of (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

If we accept the view that intraspecific competition is really predominant in population regulation, then we can understand how and why time delays destabilize stable system leading to oscillations; time delays render intraspecific competitive regulation territorially (in case of particles we take the field in which it operates) weak so that self –regulatory forces become temporarily insufficient. If we now agree that all population models should incorporate self-regulating intraspecific competition terms in surfeit strength along with other tropic intraspecific competition terms in sufficient strength along with other atrophic interactions then it is not fair to say " time delays destabilize" and we note from our analysis that it is usually a failure in self- regulation which can cause instability rather than delays by themselves.

We will give one example to illustrate our point. Consider a credit-debit system in which the debits is self –regulating by a logistic load as follows: (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

$$\frac{dx}{dt} = r\left(x - \frac{x^2}{K}\right) - axy = b_I(x) - m_I(x,y) \quad (say),$$

$$\frac{dy}{dt} = \beta xy - \delta y = b_2(x,y) - m_2(y) - \quad (say),$$

Where $K, a, \beta, \delta$ are positive constants such that the steady state $x^*\ y^*$ exists where

$$x^* = \frac{\delta}{\beta}; \quad y^* = \frac{r}{a}\left(-\frac{\delta}{K\beta}\right) > 0.$$

In this example the predator self –regulation $(\partial m_2/\partial y)$ is not sufficiently strong in the sense

$$x^* = \left(\frac{\partial m_2}{\partial y}\right)_{(x^* y^*)} \succ \left(\frac{\partial b_2}{\partial y} + \frac{\partial m_1}{\partial y}\right)_{(x^* y^*)}$$

If in such a model there is a time delay in the credit creation rate, such a delay can lead to instability of $(x^*\ y^*)$ of setting up undamped oscillations. If however we modify the growth rate by the inclusion of a density dependent self-regulation so that credit

$$\frac{dy}{dt} = \beta\, xy - \delta y - \eta y^2$$

For some positive constant η then it is possible, depending on η, for such a modified system to remain stable without oscillations even if there is a time delay in the credit creations rate. We have another completely different justification to have $\eta y^2$ in the credit dynamics: for instance if the debit is removed from this cannot happen. While this example serves to illustrate an aspect of our result, our advocacy of the significance is not based on this example.

Gopalasamy has thus derived conclusion that "time delays need not destabilize" in models with several discrete delays. However we have provided in the appendix the necessary tools (Lemmas 3,4,5) to prove all our results when the relevant delays are all (or some) continuously distributed. Some authors (Caswell [9], May [23], Cushing [12] appear to regard continuously distributed dalays as financially/cosmologically credit wise obtained when he subjected algae to a variable nitrate environment and obtained a better fit for his data with a continuous delay. The questions of whether a discrete or continuous delay is more realistic may well depend on what process is being modeled. For instance an increased supply of some resource at some point in time may trigger reproduction at a particular moment in time later irrespective of the conditions of the resource during the intervening period. If the duration of increased resource supply necessary to elicit such a response is short compared to the intervening period before the reproduction occurs, then a discrete time delay may be appropriate. In a large number of cases of animal. Note fungus algae are prey. Here it shall be an account which loses the energy, amount, or matter for that matter. Same result holds good for (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Populations of particles, atoms, molecules, propositions, systems types chosen the past dietary and nutritional history of an animal over a long period plays an influential role in determining the current bahaviour of the dietary and nutritional history of an animal over a long period plays an influential role in determining the current behavior of the animal and in such cases continuous time delays will be appropriate.

In conclusion the following comments are appropriate. The time delays since the current creation rate is usually due to a contribution of those mature animals living a few units of time earlier depending on the type of animal or organism; the time delays appearing in interspecific interactions can be due to the fact that in order to successfully interact and compete interferentially a developmental maturity may be necessary. In the case of predation, delays in the birth rate of prey can be due to delayed maturation while that in the birth rate of a predator is due to gestation delay with delayed fecundity; the delays in the predation rate are again are again as before due to the maturity needed for a predator to successfully organize and succeed in an attack.

In our analysis we have assumed throughout a temporally uniform environment neglecting dispersal of populations and their age structure. We expect our results to hold under suitable conditions both in temporally periodic environments with the steady state of a constant environment being replaced by a periodic solution

(nonstationary steady state) and when the effects of dispersion and age structures are considered. It is worthwhile to modify our purely ecological, cosmological, financial, propositional, problematic systems models so as to provide them with a genetic foundation and examine the effects of time delays since forces of selection act with some time delay. Such a generalization will definitely lead to a better understanding of "how confederations regulate themselves".

## Appendix

First we derive a comparison result.

Lemma 1. Consider the delay differential system

$$\left.\begin{aligned}\frac{dY_1(t)}{dt} &= -m_{11}\ Y_1(t) + b_{11}Y_1(t-\tau_{11}) - m_{12}\ Y_2(t-\tau_{12}),\\ \frac{dY_2(t)}{dt} &= -m_{22}Y_2(t) - m_{21}Y_1(t-\tau_{21}) + b_{22}\ Y_2(t-\tau_{22}),\end{aligned}\right\} (A.1),$$

Where the $m_{ij}$, $b_{ij}$ $(i.j = 1,2)$ are positive while $\tau_{ij}$ $(i.j = 1,2)$ nonnegative with $\tau = \max \tau_{ij}$, $i,j = 1,2$. Then

$$\left.\begin{aligned}\frac{d|Y_1(t)|}{dt} &\le -m_{11}\,|Y_1(t)| + b_{11}|Y_1(t-\tau_{11})| + m_{12}\ Y_2(t-\tau_{12})|,\\ \frac{d|Y_1(t)|}{dt} &\le -m_{22}\,|Y_2(t)| + m_{21}|Y_1(t-\tau_{21})| + b_{22}\ |Y_2(t-\tau_{21})|,\end{aligned}\right\} (A.2)$$

Where

$$\frac{d}{dt}\,|\,Y_1(t)\,| = \left\{\begin{aligned}&\frac{dY_1(t)}{dt}\ if\ \ Y_1(t) \ge 0,\\ &\frac{dY_1(t)}{dt}\ if\ \ Y_1(t) < 0,\end{aligned}\right\}\ i = 1,2. \qquad (A.3)$$

The proof is by direct verification.

LEMMA 2. Consider now (A.2) and the following system

$$\left.\begin{aligned}\frac{dZ_1(t)}{dt} &= -m_{11}\ Z_1(t) + b_{11}Z_1(t-\tau_{11}) + m_{12}\ Z_2(t-\tau_{12}),\\ \frac{dZ_2(t)}{dt} &= -m_{22}Z_2(t) + m_{21}Z_1(t-\tau_{21}) + b_{22}\ Z_2(t-\tau_{22}),\end{aligned}\right\} \qquad (A.4)$$

Together with the initial conditions

$$Z_1(s) > |Y_1(s)|,\ \ s \in [-\tau, 0],\ i = 1,2. \qquad (A.5)$$

Then

$$|Y_1(t)| < Z_1(t) \text{ for } t \ge 0 \quad \text{and} \quad i = 1,2. \qquad (A.6)$$

PROOF. Let $U_1(t) = Z_i(t) - |Y_1(s)|$, $i = 1,2.$ $t \geq -\tau.$ Then we have from (A.2) and (A.4),

$$\left.\begin{aligned} \frac{dU_1(t)}{dt} &\geq -m_{11}\ U_1(t) + b_{11}U_1(t-\tau_{11}) + m_{12}\ U_2(t-\tau_{12}), \\ \frac{dU_2(t)}{dt} &\geq -m_{22}U_2(t) + m_{21}U_1(t-\tau_{21}) + b_{22}\ U_2(t-\tau_{22}), \end{aligned}\right\} t > 0, \quad (A.7)$$

With

$$U_1(s) = Z_i(s) - |Y_i(s)| > 0, \quad i = 1,2, s \in [-\tau, 0]. \quad (A.8)$$

Consider (A.7)-(A.8) first on the interval $[\tau, 0]$ : It will follow from (A.7) that $U_1(t) \geq U_1(0) \exp [-m_{11}t]$

$$+\int_0^t [b_{11}U_1(s-\tau_{11}) + m_{12}\ U_2(s-\tau_{12})] \exp [-m_{11}(t-s)ds > 0,] \quad (A.9)$$

For all $t \in [0, \tau^*]$ since $U_i(s) > 0$ on $[-\tau^* 0]$ where $\tau^*$ = positive minimum of $(\tau_{11,}\tau_{12,}\tau_{21,}\tau_{22,})$. Similarly $U_2(t) > 0$ for $t \in [0, \tau^*]$. Now by a similar argument it will follow that $U_1(t) > 0, U_2(t) > 0$ for $t \in [\tau^*, 2\tau^*]$. Repeating the above procedure we obtain the result (A.6).

It follows from Lemmas 1 and 2 that in order to prove theorem 3.1 it is enough to prove that $Z_i(t) \to 0$ as $t \to \infty$ for $i = 1,2$ . which we will do as follows.

PROOF OF THEOREM 3.1 Define a Lyapunov-like function

$H(Z_1,Z_2)(t)=Z_1(t)+Z_2(t)$

$$+\int_{1-T}^t b_{11}\, Z_1(s)ds + \int_{1-T}^t m_{12}\, Z_2(s)ds$$

$$+\int_{1-T}^t m_{21}\, Z_1(s)ds + \int_{1-T}^t b_{22}\, Z_{22}(s)ds$$

Since each term of H in (A.10)is nonnegative for t>0, if we show that $H \to 0$ as $t \to \infty$ it will follow from Lemma 2 that

$0 \leq |\, Y_{1,}(t)\, | \leq Z_1(t) \to 0$ as $t \to \infty$,

Implying that the perturbations $Y_{1,}$, $Y_2$, governed by (A.1)approach zero as $t \to \infty$, and this will prove our result Calculate the rate of change of H along the solutions of (A.4)so that

$$\frac{dH}{dt} = \frac{\partial H}{\partial t} + \frac{\partial H}{\partial Z_1}\frac{dZ_1}{dt} + \frac{\partial H}{\partial Z_2}\frac{dZ_2}{dt}$$

$$= -[Z_1(t)\{m_{11} - (b_{11} + m_{12})\} + Z_2(t)\{m_{22} - (b_{22}m_{12})\}] < 0$$

On using (2.10)since $Z_1(t),Z_2(t)$ are positive for all t>0. Thus $H(Z_1, Z_2)(t) \to 0$ as $t \to \infty$ since H is bounded below by zero: it follows from Lemma 6 below that $Z_1(t) \to 0, Z_2(t) \to 0$ as $t \to \infty$ the proof is complete.

The following set of results are intended to be a guide to prove result similar to the above in the case of continuously distributed delays An analogue of Lemma I is the following.

LEMMA 3 Let $K_{ij}(s)$,s∈[0,∞), $i,j = 1,2$,denote nonnegative continuous functions integrable on [0,∞) normalised such that

$\int_0^\infty sk_{ij}(s)ds< \infty$; $\int_0^\infty k_{ij}(s)ds=1$, I,j=1,2.

Consider the intergrodifferential system

$$\frac{dY_1(t)}{dt} = -m_{11}Y_1(t) + b_{11}\int_0^\infty k_{ij}(s)\, Y_1(t-s)ds$$

$$-m_{12}\int_0^\infty k_{12}(s)\, Y_2(t-s)ds$$

$$\frac{dY_2(t)}{dt} = -m_{22}Y_2(t) + m_{21}\int_0^\infty k_{21}(s)\, Y_1(t-s)ds$$

$$+b_{22}\int_0^\infty k_{22}(s)\, Y_2(t-s)ds$$

*Where* $m_{ij}, b_{ij}$ $(i,j = 1,2)$ *are positive constants Then*

$$\frac{d|Y_1(t)|}{dt} \leq -m_{11}|\,Y_1(t)\,| + b_{11}\int_0^\infty k_{11}(s)|Y_1(t-s)|ds,$$

$$+m_{12}\int_0^\infty k_{12}(s)|\,Y_2(t-s)|ds,$$

$$\frac{d|Y_2(t)|}{dt} \leq -m_{22}|\,Y_2(t)\,| + m_{21}\int_0^\infty k_{21}(s)|Y_1(t-s)|ds,$$

$$+b_{22}\int_0^\infty k_{22}(s)|\,Y_2(t-s)|ds,$$

*Where* $d|Y_1(t)||/dt$ *is defined by* (A.3).

Proof is again by direct verification.

The proof of the following lemma is similar to that of Lemma 2 above.

Lemma 4. *Consider (A.14) and the following intergrodifferential system:*

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1(t) + b_{11}\int_0^\infty k_{11}(s)Z_1(t-s)ds$$

$$+m_{12}\int_0^\infty k_{12}(s)Z_2(t-s)ds,$$

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1(t) + b_{11}\int_0^\infty k_{11}(s)Z_1(t-s)ds$$

$$+b_{22}\int_0^\infty k_{22}(s)Z_2(t-s)ds,$$

Together with the initial conditions

$$\infty > Z_i(s) > |Y_i(s)| for\ s \in (-\infty, 0).$$

(A.16)

Then $|Y_i(t)| < Z_i(t) for\ t \ge 0, i = 1,2.$

The following lemma will be useful to prove a result analogous to that of Theorem 3.1 corresponding to a class of continuously distributed time delays.

LEMMA 5. Let $Z_1(t), Z_2(t)$ be the solutions of (A.15)- (A.16) where $(Y_1\ Y_2)$ corresponds to an arbitrary but fixed solutions of (A.13). Then as a consequence of Lemma & it will follow that

$$0 \le |Y_i(t)| \ Z_i(t) \to \infty \ as\ t \to \infty, \qquad i = 1,2,$$

(A.17)

Whenever

$$\begin{cases} M_{11} > b_{11} + m_{21}, \\ M_{22} > b_{22} + m_{12} \end{cases}$$

PROOF. Define a Lyapunov –like function H where

$$H(Z_1,Z_2)(t) = Z_1(t) + Z_2(t) + b_{11}\int_0^\infty k_{11}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$$

$$+ m_{12}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$$

$$+ m_{21}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$$

$$+ b_{22}\int_0^\infty k_{22}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$$

Where $(Z_1,Z_2)$is any solution of (A.15)-(A.16) for arbitrary $(Y_1,Y_2)$defined on $(-\infty, 0]$ Computing the rate of change of H in (A.19)along the solution of (A.15)we derive that (on using (A.12))

$$\left(\frac{dH(Z_1Z_2(t)}{dt} = [Z_1(t)\{m_{11} - (b_{11} + m_{21})\} + Z_2(t)\{m_{22} - (b_{22} + m_{12})\}]\right.$$

$< 0$(using (A.18)) for $Z_1(t) > 0\ and\ Z_2(t) > 0.$

By Lemma 6 below the result will follow

Lemma 6. Let H(t)=H$(Z_1, Z_2)(t)$ be as defined in (A.10) Suppose

$$m_{11} > b_{11} + m_{21} \text{ and } m_{22} > b_{22} | m_{12}$$

Then

$$Z_1(t) + Z_2(t) \to 0 \text{ast} \to \infty.$$

PROOF. Suppose (A.21) dose not hold; then there exists at least one sequence

$\{t_n\} \to \infty$ as n $\to \infty$ such that

$$[m_{11} - (b_{11} + m_{21})]\, Z_1(t_n) + [m_{22} - (b_{22} + m_{12})]Z_2(t_n) > \varepsilon \qquad (A.22)$$

For some sufficiently small positive number $\varepsilon$ and hence

$$\frac{dH(t_n)}{dt} < -\varepsilon, n = 1,2,3, \ldots.$$

By the continuity and differentiability of H with respect to t,

$$\mathrm{H}(t_n - \varepsilon) - \mathrm{H}(t_n) = \frac{dH(t^*)}{dt}(-\varepsilon), \quad t_n - \varepsilon < t^* < t_n$$

If $\varepsilon$ is small we will have

$$= \frac{dH(t^*)}{dt}\varepsilon/2$$

So that

$$\mathrm{H}(t_n) \leq \mathrm{H}(t_n - \varepsilon) - \varepsilon^2/2$$

$$\leq H(t_{n-1}) - \varepsilon^2/2$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2)$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2):$$

$$\leq H(t_0) - \mathrm{n}(\varepsilon^2/2): \mathrm{t}_0 = 0.$$

(A.12) and Lemma 5 will imply that $H(t_0) < \infty$ for any $t_0 > 0$. If n is sufficiently large it will follow that $H(t_n) < 0$ since $H(t_0) < \infty$, which contradicts the nonnegativity of H for all $t \geq 0$. Hence $Z_1(t) + Z_2(t) \to 0$ as $t \to \infty$ implying the asymptotic stability of the trivial solution of (A.1)(since $Z_1$ dominates $y_1$)

We note the result of Lemma 5 and 6 will remain valid if (A.18) is replaced by the condition $\mu(A) + \|B\| < 0$

Where A and B are defined as in (6.3).

We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities,

apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems

**Gopalsamy**

# Model For Innocuously Benign Delay In A Periodic System

**Abstract**

Sufficient conditions are obtained for the existence of a unique linearly denomination stable positive periodic solution of a system model of two competitions in a periodic environment with time lags in interspecific interactions. It is shown that if the self- regulating intraspecific interaction effects are strong enough and act without time delays then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a periodic domination. An example in finance will make things clear. Corporate finance for one big plant generates lot of accounts of the workers whose salary shall be disbursed to the debit of cash credit account. Here Cash credit account is a resource as also the one to which debits are raised, a la predator. Some of the accounts from which the credits come are demand bills purchase, usance bills discount, vostro and nostro accounts in respect of foreign bills and Letters of credit. In economics inventories are used to produce output, output is sold to produce income which might be saved in part and recycled in investment. Cosmologically take an example of a Black hole which gormandizes matter and radiates energy. Notwithstanding information paradox, one thing is clear that matter is being created in to radiation. We show how the model could be applied to axiomatic predications, postulation alcovishness, phenomenological methodologies, propositional statement, and other problems in cosmology, philosophy, and even in Green House effects. We study the sine qua non properties of such a system. In fact these models can be applied to such unsolved problems like Yang mills Theory and mass gap, and other problems. As we apply the model, we shall go on dilating the position status and the applicability of the model. We shall not touch the electrical engineering problems which might rise questions like what is charge of a Bank. We may say when a customer enters a Bank; we can say the Bank is charged. These factors are still relativistic and we shall make an in-depth study as we progress. Gopalsamy's research was published in Australian Journal of Mathematical Society vide: B 25349=365 It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 13. Introduction

At the outset we assume that the system satisfies Haimovici is caution. The purpose of this article is to establish a set of sufficient conditions for the existence of a unique linearly asymptotically stable positive periodic solution for a nonautonomous system of the form

$$\frac{dx(t)}{dt} = x(t)f_1\big(t, x(t), y(t-\tau_1)\big),$$

**(1.1)**

$$\frac{dy(t)}{dt} = y(t)f_2\,\big(t, x(t-\tau_2), y(t)\big).$$

In the context of confederation dynamics a model system of the form (1.1) means that the interspecific interactions occur with time delays $\tau_1$ 0, $\tau_2$ 0 in a periodic environment, if $f_1$ *and* $f_2$ are assumed to be periodic in $t$ with a common period $\omega$ 0. Note the interactions are inter specific.

Although model systems in periodic domains have been investigated by a large number of authors (for some references see Gopalsamy (4], [5]), periodically varying environments in time delayed model systems of two or more interacting species have not been investigated. The need for Multitype system models in periodic environments and time delayed interactions are evident from the extensive literature of that type relating to single species. While it is known that periodic ecosystems can support more species than uniform environments (Gopalsamy [4]), time delays might destabilize the system leading to a possible loss of one or more species. Furthermore, periodic environments in population dynamics have some relevance to natural selection and evolution; we will consider this aspect briefly at the end.

The principal finding of the following analysis of (1.1) is that, under suitable hypotheses on $f_1$ *and* $f_2$, the system (1.1) has a linearly asymptotically stable unique positive periodic solution for all nonnegative delays $\tau_1, \tau_2$ (while our numerical experiments suggest that the periodic solution of (1.1) is globally asymptotically stable and is independent of the delay parameters $\tau_1, \tau_2$ it is not known how one can prove such a strong result). Since time delays are usually associated with possible destabilization of otherwise stable systems, we can describe the time delays appearing on (1.1) as “harmless”.

It has been recently shown by the author (Gopalsamy [6]) that if the self –regulating negative feedback in each denomination is strong enough and acts without time delays, then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a temporally uniform domain. The result of this paper is thus an extension to a periodic environment of the "harmless" nature of interspecific interaction delays in a temporally constant environment mentioned above.

## 14. Existence and uniqueness of periodic solutions

We will now formulate the problem of existence of periodic solutions of the periodic delay differential system

$$\frac{dx_1(t)}{dt} = x_1(t) f_1\big(t, x_1(t), x_2(t-\tau_{12})\big),$$

$$t > 0,$$

**(2.1)**

$$\frac{dx_2(t)}{dt} = x_2(t) f_2\big(t, x_1(t-\tau_{21}), x_2(t)\big),$$

(Where $\tau_{12}, \tau_{21}$ are nonnegative constants) in a form suitable for the application of fixed point theory of positive operators. We make the following assumptions regarding $f_1$ and $f_2$ in (2.1).

($H_1$) The partial (Frechet) derivatives of $f_1, f_2$ satisfy the following:

$$\frac{\partial f_i}{\partial x_j} < 0, \qquad i, j = 1, 2 \text{ For all } t > -\tau \text{ where } \tau = \max(\tau_{12}, \tau_{21}). \qquad \textbf{(2.2)}$$

($H_2$) $f_1(t, x_1, x_2) \equiv f_1(t + \omega, x_1, x_2)$ for all $t > -\tau, x_1 \geq 0, x_2 \geq 0$. **(2.3)**

($H_3$) There exist positive constants $b_{il}$, $b_i^u$, $a_{il}$, $a_i^u$ ($l = 1, 2$) such that

$$b_{il} - a_{il} x_i < f_i(t, x_1, x_2) < b_i^u - a_i^u x_i,$$

$$b_{il}/a_{il} < b_i^u/a_i^u, \qquad \textbf{(2.4)}$$

For $i$= 1, 2 and $(x_1, x_2) \in R_+ \times R_+$. If $\tau = 0$ then one can show that (2.1)-(2.4) will imply that the nondelayed periodic system (2.1) will have an asymptotically stable strictly positive periodic solution (for details see Gopalsamy [5]); this result will become a special case of what we obtain in the following.

We will use the following notation:

$$B_\omega^\tau = \{h: [-\tau, \infty) \to R, h \text{ is continuous}, h(t + \omega) = h(t), t \geq -\tau\},$$

$$B = B_\omega^\tau \times B_\omega^\tau,$$

$$K_1 = \{ h \in B_\omega^\tau | h(t) \geq 0, t \in [-\tau, \infty)\},$$

$$K = K_1 \times K_1,$$

$$||h|| = \left\{ \sup_{t \in [0,\omega],} | h(t)|, h \in B_\omega^\tau \right\}$$

$$||x_1, x_2|| = \max\{||x_1||, ||x_2||, (x_1, x_2) \in B\},$$

$$G_1 = \{(x_1, x_2) \in K, ||(x_1, x_2)|| \leq b_l = \min(b_{1l} / a_{1l}, b_{2l}/a_{2l})\}$$

$$G_2 = \left\{(x_1, x_2) \in K, ||(x_1, x_2)|| \geq b^u = \max\left(\frac{b_1^u}{a_1^u}, \frac{b_2^u}{a_2^u}\right)\right\},$$

$\bar{G} =$ closure of $G_2$ *in* $B$ and

$\frac{\bar{G}_2}{G_1} =$ complement of $G_1 in\ \bar{G}_2$; $G_1 \subset G_2$ by **(2.4)**

The cone $K = K_1 \times K_1$ induces a partial order in $B$ and we will write $(x_1, x_2) \geq (y_1, y_2)$ if and only if $x_1 - y_1 \in K_1$. we define an operator $N$ on $B$ as follows

$$N(x_1, x_2)(t) = \begin{matrix} N_1(t, x_1, x_2) \\ N_2(t, x_1, x_2) \end{matrix} \quad \text{for t> } 0, \qquad \textbf{(2.5)}$$

Where

$$N_1(t, x_1, x_2) = \int_0^\omega G_{11}(m, \omega, t-s) g_1\left(m, s, x_1(x) x_2(s - \tau_{12})\right) ds$$

$$+ \int_0^1 G_{12}(m, 0, t-s) g_1\left(m, s, x_1(s), x_2(s - \tau_{12})\right) ds,$$

$$N_2(t, x_1, x_2) = \int_0^\omega G_{22}(m, \omega, t-s) g_2\left(m, s, x_1(s - \tau_{21}), x_2(s)\right) ds \qquad \textbf{(2.6)}$$

$$+ \int_0^1 G_{21}(m, 0, t-s) g_2\left(m, s, x_1(s - \tau_{21}), x_2(s)\right) ds,$$

For $t > 0$, with

$$g_1\left(m, s, x_1(s), x_2(s - \tau_{12})\right) = m x_1(s) + x_1(s) f_1\left(s, x_1(s), x_2(s - \tau_{12})\right),$$

**(2.7)**

$$g_2\big(m, s, x_1(s-\tau_{21}), x_2(s)\big) = mx_2(s) + x_2(s) f_2\big(s, x_1(s-\tau_{21}) x_2, (s)\big),$$

$$G_{ll}(m, \omega, t-s) = e^{-m\omega}(1-e^{-m\omega})^{-1} e^{-m(t-s)}, \qquad i = 1, 2,$$

**(2.8)**

$$G_{ij}(m, 0, t-s) = e^{-m(t-s)}, i \neq j, i, j = 1, 2,$$

$$N(x_1, x_2)(t) = \begin{cases} N_1(n\omega + t, x_1, x_2) \\ N_2(n\omega + t, x_1, x_2) \end{cases} for\ t \in [-\tau, 0], \qquad \textbf{(2.9)}$$

*m* being a suitably selected positive constant (more will be said below) while *n* is a positive constant such that $n\omega > \tau$.

We note immediately the following from the definition of $N$: *if* $(x_1, x_2) \in \bar{G}/G_1$ then $N(x_1, x_2)(t)$ ≡ 0 and if $N(x_1, x_2)$ is periodic $\omega$ for $t$ ≡ 0 then $N(x_1, x_2)$ is periodic in *t* with period $\omega$ for $t \geq -\tau$. The existence of at least one positive periodic solution of (2.1)-(2.4) is established by means of a few lemmas and a result of Schmitt [9].

Lemma 1. *A necessary and sufficient condition for* $(x_1(t), x_2(t))$, $t \geq -\tau$ *to be a periodic solution of (2.1)-(2.4) of period* $\omega$ *is that* $(x_1, x_2)$ *be a fixed point of N such that*

$$N(x_1, x_2)(t) = \big(x_1(t), x_2(t)\big) \quad and\ (x_1, x_2) \in K \cap \left(\frac{\bar{G}}{G_1}\right).$$

PROOF. The details of proof are similar to those of Lemma 2 of Busenberg and Cooke [1]. We first verify the continuity of $N(x_1, x_2)$ assuming that $N(x_1, x_2)(t)$ is periodic in *t* of period $\omega$. The continuity of $N(x_1, x_2)$ for $t > 0$ follows from the definition in (2.5)-(2.8) WHICH ALSO PROVIDES $\lim_{t\to 0+} N(x_1, x_2)(t)$. It will follow from (2.9) and the assumed periodicity of $N(x_1, x_2)$ *that* $N(x_1, x_2)$ is continuous

For $t \geq -\tau$. Now we will verify the periodicity of $N(x_1, x_2)$ in $t$. A simple computation leads to

$$N_1(t+w, x_1, x_2) - N_1(t, x_1, x_2) = e^{-mt} F_1(t),$$

**(2.10)**

Where

$$F_1(t) = \int_0^{\omega} G_{11}\,(m, \omega, t-s\,)g_1\big(m, s, x_1(s), x_2(s-\tau_{12})\big)ds$$

$$+\int_0^{\omega} G_{12}(m, 0, t-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$+\int_{\omega}^{\omega+t} G_{12}(m, 0, \omega-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$-\int_0^{\omega} G_{11}(m, \omega-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$+\int_0^{t} G_{12}(m, 0, -s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

**(2.11)**

Which on simplification using the periodicity of $f_1$ shows that $F_1(t) \equiv 0$ whenever $x_1, x_2, \in B$, thus implying the periodicity of $N_1$ in $t$ of period$\omega$. A similar argument for $N_2$ proves that $(x_1, x_2) \in B \Rightarrow N\,(x_1, x_2) \in B$.

Now suppose $(x_1, x_2)$ is a fixed point of *N and let* $(x_1, x_2)\ \in K \cap (\bar{G}/G_1\,)$. Then it is easily seen that $(x_1, x_2)$ is a solution of (2.1) and since $(x_1, x_2)\ \in K \cap (\bar{G}/G_1\,)$ it will follow that $(x_1, x_2)$ is a periodic solution of (2.1).

Conversely, let $(x_1, x_2)$ be a periodic solution of period $\omega$ of (2.1). Then $(x_1, x_2)$ is necessarily a periodic solution of

$$\frac{dx_1(t)}{dt} + mx_1(t) = g_1\,\big(m, t, x_1\,(t), x_2(t-\tau_{12})\big)$$

**(2.12)**

$$\frac{dx_2(t)}{dt} + mx_2(t) = g_2\,\big(m, t, x_1\,(t-\tau_{21}), x_2(t)\big),$$

For any arbitrary real constant *m* and hence for any fixed constant $m > 0$. It follows from (2.12) that any periodic solution of (2.1) and hence of (2.12) is a fixed point of $N(x_1, x_2)$. We claim that such a fixed point is in $K \cap (\bar{G}_2/G_1)$. Suppose $(x_1, x_2) \notin K \cap (\bar{G}_2/G_1)$. the following possibilities can arise:

(xiii) Either $x_1, x_2$ or both $x_1$ *and* $x_2$ may become negative for some value of *t* and it will follow from (2.1) that subsequent value of $x_1, x_2$ or $x_1$ *and* $x_2$ will be negative and hence from (2.4) and (2.1) their derivatives will be negative for all subsequent *t* , implying that $x_1, x_2$ or $x_1$ *and* $x_2$ cannot be periodic.

(xiv) While $(x_1, x_2) \in K \subset B$ there may be a value of *t*, say $t^*$, such that $x_1(t^*) > b_1^u/a_1^u$ or $x_2(t^*) > b_1^u/a_1^u$ and in such a case $x_1(t) or\ x_2(t)$ cannot be periodic since respectively $dx_1(t^*)/dt < 0\ or\ dx_2(t^*)/dt < 0,$ which will imply that the values of $x_1(t^*)\ or\ x_2(t^*)$ cannot be repeated from below.

(xv) Similarly if there is a $t^*$ such that $0 < x_1(t^*) < b_{1l}/a_{1l}\ or\ 0 < x_2(t^*) < b_{2l}/a_{2l}$ then such solutions cannot be periodic since $dx_1(t^*)dt > 0$ or $dx_2(t^*)dt > 0$, which will mean that the values of $x_1(t^*)$ or $x_2(t^*)$ cannot be repeted from above.

Thus every periodic solution of (2.1)-(2.4) is a fixed point of *N* in $K \cap \left(\frac{G_2}{G_1}\right) \subset B$.

**LEMMA 2.** *Let N be as in (2.5)-(2.6). Suppose* $\lambda \in (0,1) and\ \lambda N(x_1, x_2) = (x_1, x_2)$ *for some* $((x_1, x_2) \in K - \{0\}$. *then there exist* $\alpha_1^* > 0, \alpha_2^* > 0$ *such that*

$$0 < x_1(t) \leq \alpha_1^* < b_1^u / a_1^u,$$

$$0 < x_2(t) \leq \alpha_2^* < b_2^u / a_2^u,$$

*for* $t \in [-\tau, \infty)$.

PROOF. If $\lambda N(x_1, x_2) = (x_1, x_2)$ then from the definition of N we get for all $t > 0$ and $(x_1, x_2) \in K - \{0\}$,

$$\frac{dx_1(t)}{dt} = \lambda\, x_1\,(t) f_2(t, x_1(t), x_2\,(t-\tau_{12})) + x_1\,(t) m(\lambda - 1)$$

**(2.13)**

$$\frac{dx_2(t)}{dt} = \lambda\, x_2\,(t) f_2(t, x_1(t-\tau_{21}), x_2\,(t)) + x_2\,(t) m(\lambda - 1)$$

For all $t$ such that $x_1(t) = b_1^u/a_1^u$ or $x_2(t) = b_2^u/a_2^u$ we have $\frac{dx_1}{dt} < 0$ or $\frac{dx_2}{dt} < 0$; but these are impossible since $x_1, x_2$ are periodic in t. Thus by the continuity and periodicity of $x_1, x_2$ we have $x_1(t) \leq x_1^* < b_1^u/a_1^u$, $x_2(t) \leq x_2^* < b_2^u/a_2^u$ for all $t \geq 0$; by the definition and periodicity the same will follow for all $t \geq -\tau$. A similar argument holds when $x_1(t) < 0$ or $x_2(t) < 0$, noting that in such a case $dx_1/dt > 0$ or $dx_2/dt > 0$, again contradicting the periodicity of $(x_1, x_2)$ in *t*.

LEMMA 3. *The operator N maps* $K \cap (G_2/G_1)$ *in to K and is completely continuous*

PROOF. From the definition of N and (2.4) we derive that

$$N_1(t, x_1, x_2) = \int_0^\omega G_{11}\,(m, \omega, t-s) x_1(s)[m + b_{1l} - a_{1l} x_1(s)] ds$$

$$+ \int_0^\omega G_{12}\,(m, 0, t-s) x_1(s)[m + b_{1l} - a_{1l} x_1(s)] ds.$$

**(2.14)**

For $(x_1, x_2) \in K \cap (G_2/G_1)$, we have $0 < x_1 < b_1^u/a_1^u$ and hence we get $N_1(t, x_1, x_2) > 0$ for all $t \geq 0$ by a proper choice of the positive constant say $m_1$.

We make a similar choice say $m_2$ for *m* for the case $N_1(t, x_1, x_2) \geq 0$ and then choose *m* = max $(m_1, m_2)$. Thus we have

$$N: K \cap (G_2/G_1) \to K \qquad for\ t \geq 0,$$

And by the periodicity of N we have the same for $t \in -\tau$.

To prove the complete continuity of N we proceed as follows. We know that if $(x_1, x_2) \in B$ then $N(x_1, x_2) \in B$ and so we can restrict $t$ to the set $[0, n\,\omega], n\omega > r$. We write $N = L \circ Q$ where $L$ is the linear operator, $L : B \rightarrow B$, $L(u_1, u_2)(t)$

$$= \int_0^{\omega} G_{11}(m, \omega, t-s) \begin{bmatrix} u_1 & (s) \\ u_2 & (s) \end{bmatrix} ds, + \int_0^{t} G_{12}(m, 0, t-s) \begin{bmatrix} u_1 & (s) \\ u_2 & (s) \end{bmatrix} ds$$

**(2.15)**

And Q is the nonlinear operator

$$Q(u_1, u_2)(t) = \begin{matrix} g_1(m, t, u_1(t), u(t-\tau_{12})) \\ g_2(m, t, u_1(t-\tau_{21}), u_2(t)) \end{matrix} for\ \ t > 0,$$

**(2.16)**

$$Q(u_1, u_2)(t) = Q\,(u_1, u_2)(n\omega + t) \quad for\ t \in [-\tau, 0].$$

It can be found that

$$||L(u_1, u_2)(t)|| \leq max \left\{ \frac{m+b_1^u}{m}, \frac{m+b_2^u}{m} \right\} ||u_1, u_2||,$$

**(2.17)**

$$\left\| \frac{d}{dt} L(u_1, u_2) \right\| \leq 2\ max \left\{ \frac{m+b_1^u}{m}, \frac{m+b_2^u}{m} \right\} ||u_1, u_2||.$$

The complete continuity of the linear map $L$ follows from (2.17). Now since Q is continuous and maps bounded sets into bounded sets it will follow that $L \circ Q$ is completely continuous.

To prove the existence of periodic solution of (2.1)-(2.4) we will use the following result of Schmitt [9].

**THEOREM (Schmitt).** *Let $H_1$ and $H_2$ be nonempty bounded open neighborhoods of the origin in a Banach space B with a positive cone K such that $H_1 \subset H_2$ and let N: $K \cap (H_2 / H_1) \to K$ be completely continuous and which satisfies the following conditions:*

***(i)*** *There exists $(z_1, z_2) \in K \subset B$, $\|(z_1, z_2)\| = 1$ and a $\mu$ satisfying*

$$\mu > sup\{\|(x_1, x_2) - N(y_1, y_2)\|, (x_1, x_2) \in K \cap H_1, (y_1, y_2) \in K \cap \partial H_2\}$$ *such that all solution* $(x_1, x_2) \in K \cap (H_2/H_1)$ *of* $(x_1, x_2) = N(x_1, x_2) + \lambda(z_1, z_2) with\ 0 < \lambda < \mu\ satisfy\ (x_1, x_2) \notin \partial H_1.$

*(j) All solution* $(x_1, x_2) \in K \cap (H_2/H_1)\ of\ (x_1, x_2) = \lambda N(x_1, x_2), 0 < \lambda < 1\ satisfy\ (x_1, x_2) \notin \partial H_2.$

*Then N has at least one fixed point in* $K \cap (H_2/H_1)$.

We can now present our existence result.

**THEOREM 1.** *Assume that* $f_1, f_2$, *of* (2.1) *satisfy* (2.2)-(2.4). *then there exists at least one positive periodic solution of* (2.1)-(2.4).

PROOF. By Lemma 1, it is enough to show that the operator $N: K \cap (G_2/G_1) \to K \subset B$ has at least one fixed point in $G_2/G_1$ and we will do this by applying the result of Schmitt [9]. First we let

$$H_2 = \{(x_1, x_2) \in B | \ ||x_1, x_2|| < \max\{b_1^u/a_1^u, b_2^u/a_2^u)\} \equiv G_2,$$

$$H_1 = \{(x_1, x_2) \in B | \ ||x_1, x_2|| < \min\{b_{1l}/a_{1l}, b_{1l}/a_{1l})\} \equiv G_1.$$

By Lemma 3, $N: K \cap (H_2/H_1) \to K$ is completely continuous. Choose $\mu = \infty$ and $(z_1, z_2) \equiv (1.1)$. Then all solution $(x_1, x_2) \in K \cap (H_2/H_1)$ of $(x_1, x_2) = N(x_1, x_2) + \lambda(z_1, z_2), 0 < \lambda < \infty$, satisfy (2.1) and hence by their periodicity $(x_1, x_2) \notin \partial H_1$. The condition (a) of Schmitt's theorem is satisfied since $\mu = \infty$ by choice and $(x_1, x_2)$ and $(z_1, z_2)$ are periodic. By Lemma 2 the condition (b) of Schmitt's theorem is satisfied. Thus there exists at least

one fixed point of N in $K \cap (H_2/H_1)$ and hence there exists at least one positive periodic solution of (2.1)-(2.4) and the proof is complete.

We remark that so far no restriction on the magnitudes of the delays $\tau_{12}, \tau_{21}$ has been imposed except that they be nonnegative.

## 15. Uniqueness of the periodic solution

Applications of fixed point theorems to prove the existence of solutions of nonlinear problems usually suffer from the drawback that the uniqueness of solutions has to be proved often by other methods. We will use the following result of Smith [11] which is a slight variation of his Theorem 3.5 (see also Busenberg and Cooke [1]).

THEOREM (Smith). *Let N be an operator on a cone K in the Banach space B and let* $K^* = \{(x_1, x_2) \in K \cap B, 0 \leq$ *x1,x2≤C∗ where C∗some positive constant is. Assume the following:*

(m) If $(x_1, x_2) \in K^* - \{(0,0)\}$ and $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

(n) N is monotonic nondecreasing on $K^*$.

(o) If $(x_1, x_2)$ is in the interior of $K^*$ and if $0 < \lambda < 1$ then there exists an $\eta > 0$ depending on $(x_1, x_2)$ and $\lambda$ such that $N(\lambda x_1, \lambda x_2) \geq \lambda(1+\eta) N(x_1, x_2)$. Then N has at most one fixed point different from $(0,0)$ in $K^* \subset B$.

The cone K= $K_1 \oplus K_1$ where $K_1 = \{h \in B_\omega^\tau : h(t) \geq 0, t \in [-\tau, \infty)\}$ is slightly too wide in that solutions of the form $\{x_1(t), 0\}$ and $\{0, x_2(t)\}$ for (2.1) are also possible. We choose a narrower cone $K^*$ in $B$ as follows: define

$$S = \{(x_1, x_2) \in B : b_{1l}/a_{1l} \|x_1\| \leq b_1^u/a_1^u,\ b_{2l}/a_{2l} \|x_2\| \leq b_2^u/a_2^u,\}, \quad \textbf{(3.1)}$$

And let

$$K^* = \{\lambda(x_1, x_2) \in B : (x_1, x_2) \in S, \lambda \geq 0\}. \quad \textbf{(3.2)}$$

Since $\hat{G}_2/G_1 \subset K^*$, it is not difficult to see that N: $\hat{G}_2/G_1 \to B$ has at least one fixed point which is a periodic solution of (2.1)-(2.4). We will prove that this solution is unique by applying the result of Smith cited above. The following lemmas are needed first.

LEMMA 4. If $(x_1, x_2) \in K^* - \{0.0\}$ and if $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

PROFF. If $(x_1, x_2) = N(x_1, x_2)$ and $(x_1, x_2) \in K^* - \{0,0\}$, then $(x_1, x_2)$ is a periodic solution of (2.1) –(2.4) and hence we have

$$(b_{1l} - a_{1l}x_1)x_1 < \frac{dx_1}{dt} < (b_1^u - a_1^u x_1)x_1, \qquad (\mathbf{3.3})$$
$$(b_{2l} - a_{2l}x_2)x_2 < \frac{dx_2}{dt} < (b_2^u - a_2^u x_2)x_2,$$

And hence any nonzero periodic solution $(x_1, x_2)$ of (2.1)-(2.4) has to be such that

$$b_{1l}/a_{1l} < x_1(t) < b_1^u - a_1^u,$$

$$b_{2l}/a_{2l} < x_2(t) < b_2^u - a_2^u,$$

For $t > -\tau$, which proves the lemma.

LEMMA 5. The mapping $N$ defined by (2.5)-(2.9) is monotone nondecreasing on $K^*$.

PROOF. Let $(x_1, x_2) \in K^*$, $(y_1, y_2) \in K^*$ such that $(x_1 - y_1, x_2 - y_2) \in K^*$ and $(x_1, x_2) \neq (y_1, y_2)$. Then $x_1 \geq y_1$, $x_2 \geq y_2$. From the definition of $N$ we derive that

$N_1(t, x_1, x_2) - N_1(t, y_1, y_2)$

$$= \int_0^{\omega} G_{11}(m, \omega, t-s)\{m[x_1(s)] + x_1(s)f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s)f_1(s, y_1(s), y_2(s-\tau_{12}))\}ds$$

$$+ \int_0^t G_{12}(m, 0, t-s)$$

$$x\{m[x_1(S) - y_1(s)] + [x_1(s)f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s)f_1(s, y_1(s), y_2(s-\tau_{12}))]\}ds \qquad \mathbf{(3.4)}$$

With a similar relation for $N_2(t, x_1, x_2) - N_2(t, y_1, y_2)$. Since $f_1$ is continuous, $f$ is bounded on $K^*$; also it will follow from $(x_1, x_2) \neq (y_1, y_2)$ that

$$0 \not\equiv x_1(s) - y_1(s) \geq 0 \text{ and } 0 \not\equiv x_2(s) - y_2(s) \geq 0;$$

A sufficiently large positive constant $m$ will make the right side of (3.4) positive.A similar argument will lead to $N_2(t, x_1, x_2) - N_2(t, y_1, y_2) > 0$ perhaps with a larger constant $m$. Finally we can choose the bigger of these two constants $m$ to yield

$$N(x_1, x_2) - N(y_1, y_2) \text{ for } (x_1 - y_1, x_2 - y_2) \in K^* - \{(0,0)\}.$$

The monotonicity of $N$ on $K^*$ follows.

LEMMA 6. Let $(x_1, x_2) \in$ interior of $K^*$ and let $\lambda$ be $a$ constant such that $0 < \lambda < 1$. Suppose $f_1, f_2$ satisfy (2.2) - (2.4). Then there exists $a$ Positive constant $\eta$ such that

$$N(\lambda x_1, \lambda x_2) \geq \lambda(1+\eta) N(x_1, x_2). \quad \textbf{(3.5)}$$

PROOF. A direct computation leads to

$$N_1(t, \lambda x_1, \lambda x_2) - \lambda N_1(t, x_1, x_2)$$

$$= \lambda \int_0^{\omega} G_{11}(m, \omega, t-s) g_1\big(s, x_1(s), x_2(s-\tau_{12})\big) H_1(s, \lambda, x_1, x_2) ds$$

$$+ \lambda \int_0^{t} G_{11}(m, 0, t-s) g_1\big(s, x_1(s), x_2(s-\tau_{12})\big) H_1(s, \lambda, x_1, x_2) ds$$

**(3.6)**

Where

$$H_1(s, \lambda, x_1, x_2) = \frac{m\, x_1(s) + x_1(s) f_1\big(s, \lambda x_1(s), \lambda x_2(s-\tau_{12})\big)}{m\, x_1(s) + x_1(s) f_1\big(s, x_1(s), x_2(s-\tau_{12})\big)} - 1. \quad (\mathbf{3.7})$$

Since $f_1$ is decreasing (see (2.2)) with respect to $x_1$ and $x_2$ there exists a positive constant say $\eta_1$ such that

$$H_1(s, \lambda, x_1, x_2) \geq \eta_1 \text{ for } 0 < \lambda < 1, s > 0, (x_1, x_2) \in K^*, \quad \textbf{(3.8)}$$

And hence

$N_1(t, \lambda x_1, \lambda x_2) \geq \lambda(1+\eta_1)N_1(t, x_1, x_2)$.

A similar argument will show that there exists an $\eta_2 > 0$ such that

$N_2(t, \lambda x_1, \lambda x_2) \geq \lambda(1+\eta_2)N_2(t, x_1, x_2)$,

And the result follows if we choose $\eta = \min(\eta_1, \eta_2)$.

Now the following uniqueness result is a direct consequence of the result of Smith [11] and the Lemmas 4, 5, 6.

THEOREM 2. *Suppose $f_1, f_2$ of (2.1) satisfy (2.2) – (2.4). Then there exists a unique $\omega$ - periodic strictly positive periodic solution of (2.1).*

## 4. Linear stability of the periodic solution

Let $\{x_1^*(t), x_2^*(t), t \geq -\tau\}$ be the strictly positive unique periodic solution of (2.1). Let $\{x_1(t), x_2(t), t \geq -\tau\}$ be any other solution of (2.1). To study the linear stability of $(x_1^*, x_2^*)$ we let

$$\left.\begin{aligned} x_1(t) &\equiv x_1^*(t)\{1 + X_1(t)\}, \\ x_2(t) &\equiv x_2^*(t)\{1 + X_2(t)\}, \end{aligned}\right\} \quad t \geq -\tau, \tag{4.1}$$

And derive that the perturbations $(x_1, x_2)$ are governed by the linear variational system (after neglecting the non linear terms in $x_1, x_2$)

$$\left.\begin{aligned} \frac{dX_1(t)}{dt} &= -a_{11}(t)\, x_1^*(t)X_1(t) - a_{12}(t)x_2^*(t-\tau_{12})\, X_2(t-\tau_{12}) \\ \frac{dX_2(t)}{dt} &= -a_{21}(t)x_1^*(t-\tau_{21})x_1(t-\tau_{21}) - a_{22}(t)x_2^*(t)x_2(t), \end{aligned}\right\} \tag{4.2}$$

For $t > 0$, where

$$\left.\begin{aligned} -a_{11}(t) &= \frac{\partial f_1}{\partial x_1} \\ -a_{12}(t) &= \frac{\partial f_1}{\partial x_2} \end{aligned}\right\} \text{Evaluated at } (t, x_1^*(t), x_2^*(t-\tau_{12})), \tag{4.3}$$

$$\left.\begin{aligned} -a_{22}(t) &= \frac{\partial f_2}{\partial x_1} \\ -a_{22}(t) &= \frac{\partial f_2}{\partial x_2} \end{aligned}\right\}$$ Evaluated at $(t, x_1^*(t-\tau_{21}), x_2^*(t))$.

To establish the linear asymptotic stability of $(x_1^*(t), x_2^*(t))$ it is enough to prove that of the trivial solution of (4.2) - (4.3) We will do this by means of a combination of a comparison differential system and a Lyapunov function. First we define

$$\frac{d|X_1(t)|}{dt} = \begin{cases} \frac{dX_1(t)}{dt} & if\ X_i(t) \geq 0, \\ -\frac{d\setminus X_1(t)}{dt} & if\ X_i(t) < 0, \end{cases} \quad i = 1,2, \qquad \textbf{(4.4)}$$

And derive from (4.2)- (4.3) the following comparison system of differential inequalities:

$$\begin{aligned} \frac{d|X_1(t)|}{dt} &\leq -b_{11}, x_1^*(t)\,|X_1(t)| + b_{12}x_2^*(t-\tau_{12})\,|X_2(t-\tau_{12})|, \\ \frac{d|X_2(t)|}{dt} &\leq -b_{21}, x_1^*(t-\tau_{21})\,|X_1(t-\tau_{12})| - b_{22}x_2^*(t)\,|X_2(t)|, \end{aligned} \qquad \textbf{(4.5)}$$

For $> 0$, where

$$-b_{11} = \max_{t\in-\tau}\{-a_{11}(t)\},$$

$$\begin{aligned} b_{12} &= |\min_{t\in-\tau}\{-a_{12}(t)\}| \\ b_{21} &= |\min_{t\in-\tau}\{-a_{21}(t)\}| \end{aligned} \qquad \textbf{(4.6)}$$

$$-b_{22} = \max_{t\in-\tau}\{-a_{22}(t)\}$$

We can now formulate the result as follows. By the periodicity of $a_{ij}$, the constants $b_{ij}$ are independent of $\tau$ (see also (2.2)).

THEOREM 3. *The unique strictly positive periodic solution* $\{x_1^*(t), x_2^*(t), t \in -\tau\}$ of $(2.1) - (2.4)$ *is linearly asymptotically stable provided*

$$|-b_{11}| > b_{21} \text{ and } |-b_{22}| > b_{12},$$

**(4.7)**

*Where the constants* $b_{ij}$ $(i, j = 1,2)$ *are as defined in* (4.6).

PROOF. Consider the scalar valued Lyapunov-like function $v(t) = V(t, X_{1,}X_{2,})$where

$$v(t) = V(t, X_{1,}X_{2,}) = |X_{1,}(t)| + |X_{2,}(t)| \int_{t-\tau_{12}}^{t} b_{12} x_2^*(u)|X_{2,}(u)|du$$

$$+\int_{t-\tau_{21}}^{t} b_{21} x_1^*(u)|X_{1,}(u)|du. \quad \textbf{(4.8)}$$

Since $x_1^*, x_2^*$ are strictly positive, it is found from (4.7)-(4.) that

$$\frac{dv}{dt} \leq -[(|b_{11}| - b_{21})\, x_1^*(t)|X_{1,}(t)| + (|b_{22}| - b_{22})\, x_2^*(t)|X_2(t)|] \quad \textbf{(4.9)}$$

For $t > 0$. It follows from (4.7)-(4.9) that $V(t, X_{1,}X_{2,}) \geq |X_{1,}(t)| + X_{2,}(t)|$ and $dv/dt \leq 0$,which will imply the stability of the trivial solution of (4.2)-(4.3). We claim that (4.9) also implies the asymptotic stability of the trivial solution of (4.2)-(4.3)in the sense that

$$\lim_{t-\infty} [\,|X_{1,}(t)| + |X_{2,}(t)|\,] = 0,$$

**(4.10)**

For all bounded initial conditions $X_{1,}(s), X_{2,}(s), s \in \{\mathcal{T}, 0\}$. $Suppose (4.10) is\ not$ true Then we can find an increasing sequence $\{t_m\} \to \infty$ as $m \to \infty$ such that

$$(|b_{11}| - b_{21})\, x_1^*(t_m)|X_{1,}(t_m)| + (|b_{22}| - b_{12} x_2^*(t_m)|X_2(t_m)| > \varepsilon,$$

M=1, 2, 3…

**(4.11)**

Using the right derivative *dv/dt* of $v$ we derive that

$$v(t_m - \varepsilon) - \mathrm{v}(t_m) \leq -\int_{t_m-\varepsilon}^{t_m} \frac{\mathrm{dv}(\mathrm{T})}{\mathrm{dT}}\,\mathrm{d}\tau.$$

**(4.12)**

If $\varepsilon$ is small enough we have from (4.9)(4.11),(4.12)that $dv(\tau)/dr < (-\varepsilon/2)$ and

$$v(t_m - \varepsilon) - \mathrm{v}(t_m) \geq \varepsilon^2/2 \text{ and hence}$$

$$v(t_m) \geq v(t_m - \varepsilon) - (\varepsilon^2/2)$$

$$\geq v(t_m - 1)\text{-}(\varepsilon^2/2)$$

$$\geq v(t_m - 2)\text{-}(\varepsilon^2/2)$$

$$\geq v(t_0)\text{-m}\,(\varepsilon^2/2)$$

Thus for sufficiently large $m, v(t_m)$becomes negative which contradicts the Nonnegative nature of $v\ for\ t \geq 0$.

this contradiction shows that $|X_{1,}(\mathrm{t})| + |X_2(\mathrm{t})| \to 0$ as $\mathrm{t} \to \infty$ and the proof is complete.

As an example let us consider the case of two species Lotka- Volterra type competition in a periodic environment for for which we have

$$f_1\big(\mathrm{t}, x_1(\mathrm{t}), x_2(\mathrm{t} - T_{12})\big)1 \quad \left\{ \frac{x_1(\mathrm{t}) + a x_2(\mathrm{t} - T_{12})}{K_1(\mathrm{t})} \right\}$$

$$f_2(\mathrm{t}, x_1(\mathrm{t} - {}_{21}) x_2(\mathrm{t})) = 1 - \left\{ \frac{\beta x_1(\mathrm{t} - T_{21}) + x_2(\mathrm{t})}{K_2(\mathrm{t})} \right\}$$

Where $\alpha, \beta$ are positive constants and $k_{1,} k_{2,}$ are strictly positive continuous periodic functions with a common period Equation (4.6)and (4.7) for this special case lead to

$$\left|\operatorname{axm}_{t>T}\{-a_{11}(t)\}\right| = \max_{t>T}(1/K_1(t)) > \min_{t>T}\left(-a_{21}(\mathrm{t})\right)\Big| = \min_{t\in}, \beta/K_2(\mathrm{t})$$

$$\max_{t\in T}\{-a_{22}(t)\}\Big| = \max_{t\in T}(1/K_2(t)) > \min_{t\in T}\left(-a_{12}(\mathrm{t})\right)\Big| = \min_{t\in} \mathfrak{a}/K_1(\mathrm{t})$$

The inequalities reduce in the absence of temporal variations in $k_{1,}k_{2,}$ to the inequality 0< $\mathfrak{a}$,$\beta$<1, which is well known for the global asymptotic stability of the positive steady state corresponding to the above example in a constant systemic domain the above example shows also that delays in interspecific interactions have no influence on the local asymptotic stability of the unique periodic solution of the periodic ecosystem A number of cases of this type in a constant environment have been established recently by the author (Gopalsamy [6]).A numerical example is illustrated at the end.

**Consequences**

It can be found that all the conclusions on the existence uniqueness and stability of a periodic solution of (2.1)-(2.4) are valid when $OT_{12}$=0 or $T_{12}$=0 or both $T_{12}$=0 and $T_{21}$=0: thus it will follow that delays in interspecific competition do not destabilize an otherwise stable coexistence (oscillatory or nonoscillatory) of two competing types of interacting particles atoms molecules transactions interaction. In forthcoming work we will extend the analysis of the system (2.1)-(2.4) to a system of n-competing species with a hypothesis similar to that of (2.4) It is however an open problem to investigate the effects of delays in intraspecific competition in a periodic milieu

Pianka ([8],pages 139-140) remarks that “In real population rate of increase competitive abilities and carrying capacities do very from individual to individual with population density in space and time Indeed temporal variation in the environment may often allow coexistence by continually altering the competitive abilities of population inhabiting it Moreover time lags are doubt lessly of some importance in real population Finally a heterogeneous environment may allow real competitors to evolve divergent resource utilization patterns and to reduce interspecific competitive inhibition

To understand the mechanism of selection pressure against competitive inhibition we first write the Malthusian parameters $\mu_1$, $\mu_2$ of fitness for the two species as follows (see MacArthur[7],page 226-230):

$$\mu_1=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_1\left(t,x_1(t),x_2(t-T_{12})\right)$$

$$\mu_2=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_2(t,x_1(t-T_{21}),x_2(t))$$

Since $\left(x_1(t),x_2(t)\right)\rightarrow\left(x_1^*(t),x_2^*(t)\right)$ for large t we can write the parameters of finess in the from

$$\mu_1=f_1\left(t,x_1^*(t),x_2^*(t-T_{12})\right)$$

$$\mu_2=f_2\left(t,x_1^*(t-T_{12}),x_2^*(t)\right)$$

In these relations we can use

$$x_2^*(t-T_{12}),x_2^*(t)-T_{12}\frac{dx_2^*(t-T_{12}+\theta_1T_{12}}{dt}$$

$$x_1^*(t-T_{12}),x_1^*(t)-T_{21}\frac{dx_1^*(t-T_{21}+\theta_2T_{21}}{dt}$$

$\theta_1,\theta_2,\in(0,1)$ and rewrite $\mu_1,\mu_2$ in the form

$$\mu_1=f_1\left(t,x_1^*(T),x_2^*(t)-T_{12}\frac{dx_2^*(t-T_{12}+\theta_1T_{12}}{dt}\right.$$

$$\mu_2=f_2\left(t,x_1^*(T),-T_{21}\frac{dx_1^*(t-T_{21}+\theta_2T_{21}}{dt},x_2^*(t)\right)$$

$\theta_1,\theta_2,\in(0,1)$

It is seen that the presence of the delays $T_{12}$ and $T_{21}$ provides additional sources of variability in the finess parameters and thereby in the selection pressure against interspecific competition Since a necessary and sufficient condition for natural selection to act is differential (variable) rates of confederation growth we find that time delays coupled with periodicity of an domain may allow real to evolve divergent resource utilization patterns and to reduce

Interspecific competitive inhibition for some remarks relating to the significance of a periodic environment to sustain the coexistence of a large number of confederation types in the absence of delays we refer to a recent article (Gopalsamy)The following example has been numerically solved and the solutions are displayed graphically for two different values of the delay parameter $\mathcal{T}$:

$\frac{dx_1(t)}{dt} x_1(t)[1 - x_1(t) + 0.4x_2(t - \mathcal{T})/k_1(t)]$

$\frac{dx_2(t)}{dt} x_2(t)[1 - \{0.8\, x_1(t - \mathcal{T}) + x_2(t)\}/k_2(t)]$

Where $k_1(t) = 20 + \sin\ (\pi t)$ and $k_2(t)$=2.5+sin (πt)

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There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. “All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ....” (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies

systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational

content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification.

## GOPALSAMY'S MODEL APPIED TO COSMOLOGICAL PROBLEMS

## GLOBAL ASYMPTOTIC STABILITY IN AN ALMOST PERIODIC LOTKA- VOLTERRA SYSTEM

**Abstract**

**Gopalsamy studied** Sufficient conditions are obtained for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of a Lotka- Volterra system with almost periodic coefficients. Research was published Australian. Math Soc. B.27 (1986) 346-360. Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of

decay rate of perturbations. Sole intention, main objective and primary concern of the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are counted in the same vein as that of predator prey. Note that the model is applicable to each and every module. Quintessentially each parameter should have been given different name such as superscript 1,2, etc., as that has been done in the foregoing and a formal concatenation done at the end. Due to spatial restrictions we have not been able to do the line of thought and manner of analysis as has been done hereinbefore. This is the most important aspect of the model. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 9. Introduction

At the outset we assume that the system. This article is concerned with the derivation of a set of "easily verifiable" sufficient conditions for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of the Lotka- Volterra system

$$\frac{dx_1(t)}{dt} = x_i(t)\left\{b_1(t) - \sum_{j=1}^{n} a_{ij}\,(t)x_j(t)\right\};$$

$$t>t_0;\ t_0\ \in (-\infty,\infty);\ i = 1,2,\dots,n, \qquad (1.1)$$

where $b_1, a_{ij}\ (i,j = 1,2,\dots.,n)$ are nonnegative almost periodic functions defined for $t \in (-\infty,\infty)$.

In question satisfies Haimovici's equation based on age wise stratification given in the foregoing. Note that Haimovici's equations can be applied to Lotka-Volterra system term by term.

Let R and $R_n$ respectively denote the set of all real numbers and the $n$ - dimensional real Euclidean space; $R_n^+$ will denote the nonnegative cone of $R_n$. Since almost- periodic functions defined on R are bounded, we define the constants $b_i^l, b_l^u, a_{ij}^l, a_{ij}^u, (i,j = 1,2,\dots,n)$ by the following:

$$\inf_{t\in R} b_l(t) = b_i^l; \qquad \sup_{t\in R} b_l(t) = b_l^u, \qquad i,j = 1,2,\dots,n.$$

$$\inf_{t\in R} a_{lj}(t) = a_{lj}^l; \qquad \sup_{t\in R} a_{lj}(t) = a_{ij}^u.$$

We shall study the almost- periodic system (1.1) with the following assumption on the coefficients of (1.1);

$$b_l^l > 0; \qquad a_{ll}^l > 0, \qquad (1.2)$$

$$b_l^l > \sum_{\substack{j=1 \\ j\neq i}}^{n} a_{lJ}^u (b_j^u / a_{iJ}^l). \qquad (1.3)$$

Since the solutions of (1.1) corresponding to nonnegative initial conditions remain nonnegative subsequently, it will follow that

$$\frac{dx_1}{dt} \leq x_l\{b_l^u - a_{ll}^l x_i\}; \qquad t > t_0, i = 1,2,\dots,n, \qquad (1.4)$$

As a consequence of which we will have

$$0 < x_l(t_0) \leq b_l^u / a_{ll}^l = x_i^u \Rightarrow x_i(t) \leq b_l^u / a_{Il}^l, \qquad i = 1,2,\dots,n; t \geq t_0. \qquad (1.5)$$

(1.5) And (1.5) together lead to

$$\frac{dx_1}{dt} \leq x_l \left\{ b_l^u \sum_{\substack{j=1 \\ j\neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right) - a_{ll}^l x_i \right\}, \qquad i = 1,2,\dots,n; \qquad (1.6)$$

Now (1.2), (1.3) and (1.6) lead to

$$x_i(t_0 \leq \left( b_i^l - \sum_{\substack{j=1 \\ J\neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right) \right) (1/a_{ll}^u) = x_i^l$$

$$\Rightarrow x_l(t) \to x_i^l \ for\ t \to t_0,\ i = 1, 2, \dots, n. \quad (1.7)$$

10. It is our aim to show that the mode together with other models stated in the foregoing together with haimovicis model is applicable to whole lot of consummative debit – credit systems in cosmology, astrophysics, finance, economics, sociology, anatomy computation. We will not however divide the system studied here into 9 components/ modules due to restrictions on space….

From the foregoing preparation we have the following:

LEMMA 1.1. *Assume that* $b_i, a_{ij} (i, j = 1, 2, \dots, n)$ *are scalar nonnegative almost periodic functions defined on* $\boldsymbol{R}$ *such that (1.2) and (1.3) hold. Then the set S defined by*

$$S = \{x = (x_1, x_2, \dots, x_n) \in \boldsymbol{R}_n | x_* \to x_l \to x^*; i = 1, 2, \dots, n\} \quad (1.8)$$

*Where*

$$x_* = \min_{1 \to n} x_l^l; \qquad x_* = \min_{1 \to n} x_i^u \quad (1.9)$$

*is invariant with respect to (1.1)*

**Appellation, cognomen, compellation, designation, denomination of the conducive, infelicitous and unfortutious** responses to environment change may be influenced by changes in available domain and milieu under consideration. An example is the following study: Authors describe the gravitational degrees of freedom of the Schwarzschild black hole by one free variable. Authors introduce an equation which we suggest to be the Schrödinger equation of the Schwarzschild black hole corresponding to this model. We solve the Schrödinger equation explicitly and obtain the mass spectrum of the black hole as such as it can be observed by an observer very far away and at rest relative to the black hole. Our equation implies that there is no singularity inside the Schwarzschild black hole, and that the black hole has a certain ground state in which its mass is non-zero. **Schrodinger equation of the Schwarzschild black hole Jarmo Mäkelä (Cambridge U E-Print: gr-qc/9602008 | PDF** Particles, axions, atoms, molecules, quarks, Majorana fermions, bosons et al in its confederation capacity and aggregation processes, type interactions and interactions between landscape dynamics. Following bears ample testimony, infallible observatory to the fact there might be changes in the quantum behaviour due to merger of black holes. Following example bears enough testimony, infallible observatory and impeccable demonstration and apodictic evidence as how some of the models could be applied to cosmological problems. This is attributable and ascribable to the fact that the classification criterion is based on the age.

**How gravity kills (e) Schrödinger's cat Theorists argue that warped space-time prevents (e) quantum superpositions of large-scale objects. Elizabeth Gibney 17 June 2015 Propositions are written in model able form and can easy be scheduled.**

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| (89) If the cat in Erwin Schrödinger's famous thought-experiment behaved according to quantum theory, it would be able to exist in multiple states at once: both dead and alive. |
| (90) Physicists' common explanation for why we don't (e) see such quantum superpositions — in cats or any other aspect of the everyday world — is interference from (e) the environment. |
| (91) As soon as a quantum object interacts with (e&eb) a stray particle or a passing field, it picks (e) just one state, collapsing (eb) into our classical, everyday view. |
| (92) But even if physicists could completely isolate (eb) a large object in (eb) a quantum superposition, according to researchers at the University of Vienna, it would still collapse into (e&eb) one state — on Earth's surface, at least. |
| (93) "Somewhere in interstellar space it could be that the cat has (e) a chance to preserve quantum coherence, but on Earth, or near any planet, there's (=) little hope of that," says Igor Pikovski. The reason, he asserts, is gravity. |
| (94) Pikovski and his colleagues' idea, laid out in a paper published in Nature Physics on 15 June1, is at present only a mathematical argument. But experimenters hope to test whether gravity really does collapse (eb) quantum superpositions, says Hendrik Ulbricht, an experimental physicist at the University of Southampton, UK. "This is a cool, new idea, and I'm up for trying to see it in experiments," he says. Assembling the technology to do so, however, may take as long as a decade, he says. |
| (95) Cinema-goers who saw the film Interstellar are already familiar with the basic principle behind the Vienna team's work. Einstein's theory of general relativity states that an extremely massive object causes clocks near it to (e) run more slowly because (e) its strong gravitational field stretches (e&eb) the fabric of space-time (which is why a character in the film aged only an hour near a black hole, while (e) seven years passed on Earth). On a subtler scale, a molecule placed nearer the Earth's surface experiences (e&eb) a slightly slower clock than one placed slightly further away. |
| (96) Because of gravity's effect on (e&eb) space-time, Pikovski's team realised that variance in a molecule's position will also influence (e&eb) its internal energy — the vibrations of particles within (eb) the molecule, which evolve (eb) over time. |
| (97) If a molecule were put in a quantum superposition of two places, the correlation between position and (e&eb) internal energy would soon cause (eb) the duality to 'decohere' to (e) the molecule taking just one path, they suggest. |
| (98) "In most situations decoherence is due to (e) something external; here it's as though the internal jiggling is interacting with (e&eb) the motion of the molecule itself," adds Pikovski. |

**A practical limit**

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| (99) No one has yet seen this effect because other sources of decoherence — such as magnetic fields, thermal radiation and vibrations — are (=) typically much stronger and cause (eb) quantum systems to collapse long before gravity becomes an issue. But experimenters are keen to try. |
| (100) Markus Arndt, an experimental physicist also at the University of Vienna, already tests whether quantum superpositions can be observed for (e) large, though not cat-sized, objects (see Fattening up Schrödinger's cat). |
| (101) He sends large molecules through (e&eb) a matter-wave interferometer, a system that gives (eb) each molecule a choice of (e) two different paths. |
| (102) In the classical view, a molecule travels down just one path; a quantum molecule effectively passes through (e&eb) both routes at once and interferes with (e&eb) itself to create (eb) a characteristic wave-like pattern (see image, below). |

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| (103) | A quantum interference pattern built up from complex molecules known as phthalocyanines.(deleted due to spatial constraints) |
| (104) | A similar set up could be used to (e) test for gravity's ability to destroy (e) quantum behaviour: by comparing a vertical interferometer, in which the superposition should soon decohere due to (e) the stretching of time in one path versus (e&eb) the other, with a horizontal set up, where (e) the superposition could remain. |
| (105) | Arndt, who has tested the effect for (e) molecules of up to 810 atoms2, points out those large molecules, would be (=) good for testing the gravitational effect because (e) they contain (e) lots of particles contributing to (e) internal energy. |
| (106) | **But not only would the researchers have to further suppress the external environment to reduce (e) other decoherence effects, they would also need to (e) increase the separation of the two paths from micrometres to metres, or else use molecules perhaps 1 million times more massive. "It's certainly very challenging," says Arndt**. |
| (107) | If gravity's effect does limit (e) quantum behaviour on Earth, tests of quantum reality for (e) large objects may eventually have to (e) move into space, says Angelo Bassi, a physicist at the University of Trieste in Italy. |
| (108) | "But from a deep, fundamental point of view, this is nothing new," he says. A gravitational field is merely another environment to interact with, so invoking it does not (e) explain whether quantum behaviour might lead to (eb) classical reality if gravity's influence were mitigated — for example, by doing the experiment in (eb) gravity-free space. |
| (109) | Nor does the effect described by Pikovski and colleagues say anything about quantum gravity: a theory that would unite gravity and (e&eb) quantum mechanics in (eb) a single description, which many researchers are working towards. |
| (110) | "It's an interesting effect, but it's still quantum physics applied to (e&eb) classical general relativity. In that way it doesn't change our picture of the world," adds Bassi. Nature doi:10.1038/nature.2015.17773 |

Integrated factors these influences because they deal with ambience, mise en scene change or confederation dynamics, but rarely both. Time series of (e) space time suitability models with spatially explicit stochastic models to explore factors that influence the viability of type viability and confederation feasibility under stable and changing scenarios. Methodology could be applied to interrogation propositions and problems which might subsume one with the other notwithstanding that there is a space of singularities and there is space for completion of present from the past and future. An example is Time image movement from Deleuze. We clearly face a stiff challenge in our attempt to proffer Deleuzian treatments of photography as our critical project brushes up against the fact that Deleuze does not invest photography with (e&eb)the creative and re-creative potential that he associates with (e&eb) both cinema and modern painting. In short, while he highly values other visual artistic forms, he seemingly presents photographic texts as stagnate documents or tools that produce certainty, organize bodies and desires, and iterate hackneyed ideas. He even uses photography as something of a foil to demonstrate the innovation of the cinema and the originality of modern painters. In Francis Bacon: The Logic of Sensation (1981), he notes how "photography has taken over the illustrative and documentary role, so that modern painting no longer needs to fulfill this function, which still burdened earlier painters" (10). He later explains that "photographs are ways of seeing, and as such, they are illustrative and narrative reproductions or representations ... . They are what are seen, until finally one sees nothing else." He treats the photo as an instrument for reproducing representations of reality—a device that iterates images until they are ossified as established stories, icons, or even stagnant perceptions. Indeed, like Bergson before him, Deleuze compares photography to what he identifies as a reductionist model of perception: "what we see, what we perceive, are photographs" (74).[4] The photo is, in effect, an always-already passé

sensational experience that numbs our sensitivity to the ongoing vitality of life and provides us with an efficient understanding of the dynamism of our world. But it is compelling that Deleuze theorizes both the quotidian nature and the liberatory potential of photography; he treats it as a common, everyday visual experience that, despite its aesthetic and technological limitations, emancipates the modern painter to explore new artistic opportunities. This special issue of Rhizomes repeatedly provides artistic examples and critical strategies that invite us to see photography as replete with the power to unburden and emancipate that Deleuze associates with both film and painting. Indeed, one way to think of this project of reading Deleuze against Deleuze might be to ask: if photography can free painting and painters to pursue novel creative opportunities, can it likewise free itself, photographic practitioners, and even viewers to create and re-create anew? **Rhizomes » Issue 23 (2012) » Michael Kramp Unburdening Life, or the Deleuzian Potential of Photography Michael Kramp Lehigh University** dmk209@lehigh.edu

**Existence of an almost- periodic solution**

DEFINITION 2.1. A function $f(t,x)$,where $f$ is an $m$- vector, $t$ is a real scalar and $x$ is an $n$- vector, is said to be almost periodic in $t$ uniformly with respect to $x \in X \subset R_n$, $if\ f(t,x)$ is continuous in $t \in R\ and\ x \in X$, and if for any $\varepsilon > 0$, it is possible to find a constant $l(\varepsilon) > 0$ such that in any interval of length $l(\varepsilon)$ there exists a $\tau$ such that the inequality

$$\left||f\ (t+\ \tau,x) - \ f(t,x)|\right| = \sum_{l=1}^{m} |f_l\ (t+\ \tau,x) = \ f_l(t,x)|$$

$$< \varepsilon \qquad (2.1)$$

is satisfied for all $t \in (-\infty,\infty), x \in X$. The number $\tau$ is called an $\varepsilon -$ translation number of $f(t,x)$.

Let $\{\lambda_j\}$ denote the set of all real numbers such that

$$\lim_{T\to\infty} T^{-1} \int_0^T f(t,x)\exp(-i\,\lambda_j t)dt \not\equiv 0 \quad for\ x \in X; i = \sqrt{-1}. \qquad (2.2)$$

It is known that when the set X is separable, the set of numbers $\{\lambda_j\}$ in (2.2) is countable. The set $\{\sum_1^N n_j\ \lambda_j\}$ for all integers N and integers $n_j$ is called the module of $f(t,x)$.

The following result concerned with module containment will be used in the proof of our existence theorem below.

**MODULE CONTAINMENT THEOREM** (Yoshizawa [7], page 18). *Let $f(t,x)$ and $g(t,x)$ be almost periodic in t uniformly for $x \in D \subset \boldsymbol{R}_n$. If for any compact set $S \subset D$ and for any sequence of real numbers $\{\tau_k\}$ having a limit as $k \to \infty$ (the limit being finite or$\pm\infty$) for which $\{f(t+\tau_k,x)\}$ is uniformly convergent on $\boldsymbol{R} \times$*

$S, \{g(t+,x)\}$ *is also uniformly convergent on* $\boldsymbol{R} \times S$ *then the module of* $g(t+,x)$ *is contained in the module of* $f(t+,x)$.

We can now formulate our existence theorem.

**THEORM 2.1** *Assume that the almost-periodic parameters* $b_1, a_{1J}$ $(i,j = 1,2,\dots,n)$ *of (1.1) satisfy (1.2)-(1.3). Furthermore assume that there exists a positive number* $\mu$ *such that*

$$\inf_{1 \in \boldsymbol{R}} a_{11}(t) > \sum_{\substack{J=1 \\ J \neq 1}}^{n} (\sup_{t \in \boldsymbol{R}} a_{Jl}(t)) + \mu, \quad i = 1,2,\dots,n. \tag{2.3}$$

*Then the system (1.1) has a strictly positive (component wise) almost-periodic solution, say* $w(t) = w_1(t), \dots, w_n(t), t \in \boldsymbol{R}$, *whose module is contained in that of*

$$f(t,x) = \{f_1(t,x_1,\dots,x_n), f_2(t,x_1,\dots,x_n), \dots, f_n(t,x_1,\dots,x_n)\}$$

*Where*

$$f_l(t,x_1,\dots,x_n) = x_l \left\{ b_l(t) - \sum_{j=1}^{n} a_{lj}(t) x_J \right\},$$

$$i = 1,2,\dots n; (x_1,x_2,\dots,x_n) \in S. \tag{2.4}$$

PROOF. The set $S$ defined by (1.8)-(1.9) is compact in $\boldsymbol{R}_n$. Let $\{\tau_h\}$ be any sequence of real numbers such that $\tau_k \to \infty$ as $k \to \infty$ and

$$\sum_{l=1}^{n} |f_i(t+\tau_k, x_1, \dots, x_n) - f_l(t, x_1, \dots, x_n)| \to 0 \quad \text{as } k \to \infty \tag{2.5}$$

Uniformly in $(x_1, \dots, x_n) \in S$ for $t \in (-\infty, \infty)$. for any real number $\beta$let $k_0(\beta)$ be the smallest value of $k$ such that $\tau k_0 + \beta \geq t_0$. Since the invariance of $S$ implies that any solution of (1.1)

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), x_2(t+\tau_k, t_0, x_2(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\} \in S$$

For all $t \geq \beta$ and $k \geq k_0(\beta)$ whenever $\{x_1(t_0), \dots, n_n(t_0)\} \in S$, our first task in the proof is to show that the sequence

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Converges to a continuous bounded function say $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined for $t \in [\beta, \infty)$ such that $w_1(t) \in [x_*, x^*], t \in [\beta, \infty), i = 1, 2, \dots, n$ and the convergence is uniform on all compact subsets of $[\beta, \infty)$. Since $x_*$ and $x^*$ are independence of $\beta$, it is enough to show that

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Is a Cauchy sequence on compact subsets of $[\beta, \infty)$.

Let $U$ be any compact subset of $\beta, \infty)$., and let be any arbitrary positive number. Choose an integer $n_0(\varepsilon, \beta) \geq k_0(\beta)$ so large that for $m \geq k \geq n_0(\varepsilon, \beta)$

and all $t \in (-\infty, \infty)$

$$\sum_{i=1}^{n} |b_1(t+\tau_m) - b_1(t+\tau_k)| < (\mu\varepsilon x_*)/(4x^*), \tag{2.6}$$

$$x^* \sum_{i=1}^{n}\sum_{j=1}^{n} |a_{lj}(t+\tau_m) - a_{lj}(t+\tau_k)| < (\mu\varepsilon x_*)/(4x^*) \tag{2.7}$$

$$(2nx^*/x_*)\exp[-\mu x_*(\beta + \tau_k - t_0)] < (\varepsilon/2x^*)$$

(2.8)

Denote $x_i(t, t_0, x_1, (t_0))$ $by$ $x_1(t)$ brevity, and consider a function $v(s)$ defined by

$$v(s) = v\left(x(s), x(s+\tau_m+\tau_k)\right)$$

$$= \sum_{i=1}^{n} |\log x_1, (s) - \log x_1(s+\tau_m+\tau_k)|, \qquad x \geq t_0, \tag{2.9}$$

Where $x(t) = \{x_1(t), \dots, x_n(t)\}$ denotes a solution of (1.1) with $\{x_1, (t_0), \dots, x_n(t)\} \in S$. the invariance of S for (1.1), together with the elementary mean value theorem of differential calculus, implies

$$(1/x^*)\sum_{l=1}^{n}|x_l(s)-x_1(s+\tau_m+\tau_k)|\ \upsilon(s)$$

$$(1/x^*)\sum_{i=1}^{n}|x_l(s)-x_1(s+\tau_m+\tau_k)|. \tag{2.10}$$

Calculating the right derivative $D^+\upsilon$ of $\upsilon$ along the solutions of (1.1) and simplifying

$$D^+\upsilon(s)=\sum_{i=1}^{n}\{sign[x_i(s)-x_i(s+\tau_m-\tau_k)]\}$$

$$\left\{\frac{d[logx_i(s)]}{ds}-\frac{d[logx_i(s+\tau_m-\tau_k)]}{ds}\right\}$$

$$\sum_{l=1}^{n}|b_l(s)-b_l(s+\tau_m-\tau_k)|$$

$$+\sum_{l=1}^{n}(-a_{il}(s)|x_i(s)-x_i(s+\tau_m-\tau_k)|+\sum_{\substack{J=1\\ j\neq l}}^{n}a_{ij}(s)|x_j(s)-x_j(s+\tau_m-\tau_k)|)$$

$$+\sum_{i=1}^{n}\sum_{j=1}^{n}|a_{ij}(s)|a_{ij}(s+\tau_m-\tau_k)|\,x_j(s+\tau_m-\tau_k) \tag{2.11}$$

$$-\mu x_*\upsilon\big(x(s),x(s+\tau_m-\tau_k)\big)+\sum_{l=1}^{n}|b_l(s)-b_i(s+\tau_m-\tau_k)|$$

$$+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|. \tag{2.12}$$

Let us set $s-\tau_k=\sigma$; then we have from (2.6)-(2.7) that

$$\sum_{i=1}^{n}|b_l(s)-b_1(s+\tau_m-\tau_k)|+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|.$$

$$= \sum_{l=1}^{n} |b_l(\sigma + \tau_k) - b_i(\sigma + \tau_m)|$$

$$+x^* \sum_{i=1}^{n} \sum_{J=1}^{n} |a_{ij}(\sigma + \tau_k) - a_{ij}(\sigma + \tau_m)|$$

$$< (\mu\varepsilon x^*)/(2x^*). \quad (2.13)$$

Now an integration of (2.12) over $(t_0, t - \tau_k)$ together with (2.13) leads to

$$v\big(x(t+\tau_k), x(t+\tau_m)\big) \leq v\big(x(t_0), x(t_0 + \tau_m - \tau_k)\big) \exp[-\mu x_*(t + t_k - t_0)]$$

$$+[(\mu\varepsilon x_*)/(2x^*)]/\mu x_*)$$

$$\leq [(2x_n^*)/x_*]\exp[-\mu x_*(t + t_k - t_0)] + (\varepsilon/2x^*)$$

$$< (\varepsilon/x^*) \text{ On using (2.8), (2.10)} \quad (2.14)$$

For $m \geq n_0(\varepsilon, \beta)$ and for all $t \in \cup \subset [\beta, \infty]$, Thus we have from the definition of $v$ and (2.14) that

$$\sum_{l=1}^{n} |log x_i(t + \tau_k) - log x_l(t + \tau_m)| < (\varepsilon/x^*) \quad (2.15)$$

And hence using $x_* \leq x_i(t) \leq x^*, i = 1, 2, \dots, n; t \geq t_0$,

$$\sum_{i=1}^{n} |x_i(t + \tau_k) - x_i(t + \tau_m)| < \varepsilon \text{ for all } t \in \cup \subset [\beta, \infty], m \geq n_0(\varepsilon, \beta). \quad (2.16)$$

The existence of a function $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined on $[\beta, \infty]$, such that $w_i(t) \in [x_*, x^*], i = 1, 2, \dots, n; t \in [\beta, \infty]$ follows from (2.16) and the fact that $x_i(t) \in [x_*, x^*], i = 1, 2, \dots, n;$ for $t \geq t_0$. Since $\beta$ is arbitrary, $w$ is in fact

Defined on $(-\infty, \infty)$ and hence we have

$x_i(t+\tau_k) \to w_l(t)$ as $k \to \infty$ uniformly on compact subsets of $(-\infty,\infty)$, $i = 1,2,\dots,n.$

(2.17)

Our next task is to show that $w: R \to S$ is differentiable and is a solution of (1.1). Since $x(t) = \{x_1(t),\dots,x_n(t)\}$ satisfies (1.1) we have $\sum_{i=1}^{n} |\frac{d}{dt} x_i(t+\tau_k) - \frac{d}{dt} x_i(t+\tau_m)|$

$$\sum_{i=1}^{n} |f_l\left(t+\tau_k, x_1(t+\tau_k),\dots,x_n\left(t+\tau_k\right)\right)$$

$$- f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)| \qquad (2.18)$$

$$+\sum_{l=1}^{n} |f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$- f_l\left(t+\tau_m, x_1(t+\tau_m),\dots,x_n(t+\tau_m)\right)|$$

Where $f_l\ (i = 1,2,\dots,n)$are as in (2.4). Since $x(t+\tau_k) = x_1(t+\tau_k),\dots, x_n(t+\tau_k) \in S$ for large $\tau_k$ for $t$ in each compact subset of $(-\infty,\infty)$ there exists an integer $n_1(\varepsilon) > 0$ such that if $m \text{ } k \text{ } n_1(\varepsilon)$, then for any $\varepsilon > 0$,

$$\sum_{i=1}^{n} |f_l\left(t+\tau_k, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$-f_1\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)| < (\varepsilon/2). \qquad (2.19)$$

Similarly since $x_l(t+\tau_k) \in S$ and $x_l(t+\tau_m) \in S$ for $i = 1,2,\dots,n$ there is an integer $n_2(\varepsilon) > 0$ such that for large enough $m, k$, such that $m \text{ } k \text{ } n_2(\varepsilon)$, we have

$$\sum_{l=1}^{n} |f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$-f_1\left(t+\tau_m, x(t+\tau_m),\dots,x_n(t+\tau_m)\right)| < (\varepsilon/2) \qquad (2.20)$$

This equation will be very much useful in chemistry if used with the modules mentioned and written in concatenation with each other. It should be borne in mind that the applicability is modulewise and in no case concatenation should be marginalized. In all cases parameters are different which could be represented with different superscripts and subscripts like 1,23,…….etc., **Anal Bioanal Chem. 2004 Sep; 380(1):36-45. Epub 2004 Jul 29. Electronic tongue for pharmaceutical analytics: quantification of tastes and masking effects Legin A et al**

The organoleptic aspects of pharmaceutical formulations affect their acceptability to the patient and hence can have an important effect on concordance with treatment. Objective evaluation of these aspects, particularly the taste of the formulation and the drug substance it contains, is difficult. Whilst volunteer taste panels can be used to good effect their utility is limited, particularly during very early stage development when the toxicological profile of the active pharmaceutical ingredient (API) is yet to be established in detail. A potentiometric "electronic tongue" has been applied to analyse a variety of 41 individual substances and mixtures of particular interest for pharmaceutical research and development. The electronic tongue (ET) was capable of discriminating between substances with different taste modalities and could also distinguish different substances eliciting the same basic taste; the ET is promising in terms of quantifying the content of each substance and has an ability to detect nuances of the basic taste (e.g. lingering or short-lived). After calibration the electronic tongue was successfully applied to predicting bitterness strength of binary mixtures with a sweetener in terms of "apparent" or "perceived" quinine content. In order to render a formulation palatable it is often necessary to mask the (usually bitter) taste of the API by the addition of masking agents such as sweeteners and flavours. The ET proved capable of distinguishing between formulations with different levels of sweetener and/or flavour in a manner that was consistent with their masking efficiency as perceived by a small human taste panel. A suitably calibrated ET could have the benefit of providing the pharmaceutical formulator with reliable data concerning the taste of the product quickly and with a reduced need to ask volunteers to taste active pharmaceutical samples. Early development activities could be facilitated when human tasting is usually not possible in the absence of the required toxicological data.

Thus if $m \; k \; n_0(\varepsilon) = max\{n_1, n_2\}$, we have

$$\sum_{l=1}^{n} |\frac{d}{dt}\, x_1(t+\ \tau_k) - \frac{d}{dt} x_1\ (\ (t+\tau_m)| < \varepsilon \qquad (2.21)$$

Which shows that

$$\lim_{k\to\infty} \frac{d}{dt} x_1(t+\ \tau_k) \quad exists\ for\ i = 1, 2, \dots, n$$

Uniformly on all compact subsets of $(-\infty, \infty)$. Thus we have

$$\lim_{k\to\infty} \lim_{k\to\infty} \frac{x_i(t+\tau_k+h) - \ x_i(t+\tau_k)}{h}$$

$$= \lim_{h\to 0} \lim_{k\to 0} \frac{x_l(t+\tau_k+h) - \ x_l(t+\tau_k)}{h} \qquad (2.22)$$

$$= \lim_{h\to 0} \frac{w_1(t+h) - \ w_1(t)}{h}, \qquad i = 1, 2, \dots, n,$$

Showing that (since the left side of (2.22) exists) $w_i (i = 1, 2, \dots, n)$ is differentiable. Furthermore we have from (2.22) that

$$\frac{dw_l(t)}{dt} = \lim_{k\to\infty} [\, f_i\ \big(t+\tau_k, x_1(t+\tau_k), \dots, x_n\ (t+\tau_k)\big)$$

$$-f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big)$$

$$+ - f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big)]$$

$$= \ -f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big); \qquad i = 1, 2, \dots, n, \qquad (2.23)$$

By the uniform almost periodicity of $f_1, f_2, \dots, f_n$ on $R \times S$ and hence $w(t) = \big(w_1(t), \dots, w_n(t)\big)$is a solution of (1.1).

In order to complete the proof, we have to show that $w(t) = \big(w_1(t), \dots, w_n(t)\big)$ is an almost- periodic function whose module is contained in that of $\{f_1(t, x_1, \dots, x_n), f_2(t, x_1, \dots, x_n), \dots, f_n(t, x_1, \dots, x_n)\}$ for $(x_1, x_2 \dots, x_n) \in S$. For this purpose it is in fact sufficient (by the module containment theorem) to show that for any sequence of real numbers $\{\tau_k\}$ for which

$$\{f_1(t + \tau_k, x_1, \dots, x_n), f_2(t, + \tau_k, x_1, \dots, x_n), \dots, f_n(t, + \tau_k, x_1, \dots, x_n)\}$$

Converges uniformly on R for all $(x_1, \dots, x_n) \in S$, the sequence $\{w_1(t, + \tau_k) \dots, w_n(t + \tau_k)\}$ converges uniformly on R where $\tau_k$ tends to a finite number or $\pm\infty$ as $k \to \infty$.

In the case where $\tau_k \to \tau \neq \pm\infty$ as $k \to \infty$, we have $w_l(t, + \tau_k) \to w_i(t + \tau), i = 1, 2, \dots, n$ uniformly on R as $k \to \infty$. Let us suppose $\tau_k \to \infty$ as $k \to \infty$. For any $\varepsilon > 0$ there exists an integer $n_0(\varepsilon)$ such that if $m$ [illegible] $n_0(\varepsilon)$ then

$$\sum_{l=1}^{n} |b_l(t + \tau_k) - b_l(t + \tau_m)| < (\varepsilon/4x^*)x_*, \qquad (2.24)$$

$$x^* \sum_{i=1}^{n} \sum_{j=1}^{n} |a_{ij}(t + \tau_k) - a_{ij}(t + \tau_m)| < \left(\frac{\varepsilon}{4x^*}\right)\mu x_* \qquad (2.25)$$

$$n(2x^*/x_*) \exp[\mu x_* \tau_k] < (\varepsilon/2x^*). \qquad (2.26)$$

Consider the function $v\big(w(s), w(s + \tau_m - \tau_k)\big)$ for $s \in [t, t + \tau_k]$ defined by

$$v\big(w(s), w(s + \tau_m - \tau_k)\big) = \sum_{l=1}^{n} |\log w_l(s) - \log w_l(s + \tau_m - \tau_k)|. \qquad (2.27)$$

Using the fact $w: \boldsymbol{R} \to S$ and $w$ is a solution of (1.1), calculate the right derivative $D^+v \ of \ v$;

$$D^+v\big(w(s), w(s + \tau_m - \tau_k)\big) = \sum_{i=1}^{n} \{\text{sign}[w_l(s) - w_i(s + \tau_m - \tau_k)]\}.$$

$$\left\{\frac{dw_i(s)}{ds} \cdot \frac{1}{w_l(s)} - \frac{dw_l(s + \tau_m - \tau_k)}{ds} \frac{1}{w_l(s + \tau_m - \tau_k)}\right\}$$

$$\leq -\mu \sum_{l=1}^{n} |w_l(s) - w_l(s+\tau_m-\tau_k)| + \sum_{l=1}^{n} |b_i(s) - b_l(s+\tau_m-\tau_k)|$$

$$+x^* \sum_{l=1}^{n}\sum_{J=1}^{n} |a_{lJ}(s) - a_{lJ}(s+\tau_m-\tau_k)| \quad (2.28)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+ \sum_{i=1}^{n} |b_l(s) - b_l(s+\tau_m-\tau_k)| \quad (2.29)$$

$$+x^* \sum_{l=1}^{n}\sum_{j=1}^{n} |a_{lJ}(s) - a_{ij}(s+\tau_m-\tau_k)|.$$

If we set $s-\tau_k = \sigma$ then we have from (2.29),

$$D^+ v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+ \sum_{l=1}^{n} |b_l(\sigma+\tau_k) - b_l(s+\tau_m)|$$

$$+x^* \sum_{l=1}^{n}\sum_{j=1}^{n} |a_{ij}(\sigma+\tau_k) - a_{ij}(\sigma+\tau_m)| \quad (2.30)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+[(\varepsilon\mu x_*)/(4x^*)] + [(\varepsilon\mu x_*)/(4x^*)]$$

(2.31)

Integrating (2.31) over $(t, t+T_k)$

$$v(w(t+T_k)w(t+T_m))$$

Another MODEL where the model is directly applicable with the scheduled ones are the case of insurance and corporate finance. Insurance: Mathematics and Economics Volume 42, Issue 1, February 2008, Retrieval of Black–Scholes and generalized Erlang models by perturbed observations at a fixed time Dedicated to Professor André Dubey on the occasion of his 60th birthday Daniel Neuenschwander

S-stable laws on the real line (more generally on Hilbert spaces), associated with some non-linear transformations (so-called "shrinking operations"), were introduced in [Jurek, Z.J., 1977. Limit distributions for truncated random variables. In: Proc. 2nd Vilnius Conference on Probability and Statistics, June 28–July 3, 1977. In: Abstracts of Communications, vol. 3, pp. 95–96; Jurek, Z.J., 1979. Properties of s-stable distribution functions. Bull. Acad. Polon. Sci. Sér. Math. XXVII (1), 135–141; Jurek, Z.J., 1981. Limit distributions for sums of shrunken random variables. Dissertationes Math. vol. CLXXXV]. In [Jurek, Z.J., Neuenschwander, D., 1999. S-stable laws in insurance and finance and generalization to nilpotent Lie groups. J. Theoret. Probab. 12 (4), 1089–1107], the authors interpreted s-stable motions on the real line as limits of total amount of claims processes (up to a deterministic premium) of a portfolio of excess-of-loss reinsurance contracts and showed that they led to Erlang's model or to Brownian motion. In [Neuenschwander, D., 2000b. On option pricing in models driven by iterated integrals of Brownian motion. In: Mitt. SAV 2000, Heft 1, pp. 35–39], we considered stochastic integrals whose integrand and integrator are both independent Brownian motions, thus modelling a stochastic volatility; as a result we got an analogue of the Black–Scholes formula in this model, confirming a result of Hull and White [Hull, J., White, A., 1987. The pricing of options on assets with stochastic volatility. J. Finance XLII (2), 281–300]. In the present paper, we will look at a common generalization of these processes, namely s-stable motions on the real line perturbed by a stochastic integral whose integrand and integrator are both (not necessarily independent) s-stable motions. The main result will be that if we can observe the distribution of such so-perturbed s-stable motions (together with the values of the perturbing processes) at time t=1, then we can identify the whole model (including the perturbation) among all models with Lévy processes perturbed by an iterated stochastic integral of two Lévy processes (in the Gaussian case) resp. among all models with a compound Poisson process with drift perturbed by an iterated stochastic integral of two compound Poisson processes (in the completely non-Gaussian case if the perturbing processes have no drift) without knowing anything about the history or about its distribution during $0 \le t < 1$. This applies, e.g., to a situation

where several assets obey the same model and one can estimate the distribution at time one by looking at the values of all these assets at time t=1. Interestingly enough, it will be convenient to treat the whole matter in the algebraic framework of the so-called Heisenberg group. This is a concept coming in fact from quantum mechanics and is in a certain sense the simplest non-commutative Lie group.

$\leq e^{-\mu x^* T_k} v\big(w(t), w(t+T_m-T_k)\big) + (\varepsilon/2x^*)$

$\leq e^{-\mu x^* T_k} \sum_{i=1}^{n} |\log\ w_i(t) - \log\ w_i(t+T_m-T_k)| + (1/x_*) + (\varepsilon/2x^*)$

$\leq n(2x^*/x_*)\exp[-\mu x_* T_k] + (\varepsilon/2x^*)$

$<(\varepsilon/x^*)$.

Since

$$v\big(w(t+T_k)w(t+T_m)\big) \leq \sum_{i=1}^{n} |\ w_1(t+T_k) - w_1(t+T_m))/x^*)$$

We have from(2.33)that

$$\sum_{i=1}^{n} |\ w_1(t+T_k) - w_1(t+T_m)| < \varepsilon\ for\ m \geq k \geq n_0(\varepsilon)$$

Which show that the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\}$converges uniformly on **R**. By a similar argument the above conclusion is also valid when $T_k \to -\infty\ as\ k \to \infty$. Thus the uniform convergence of the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\}\ as\ k \to \infty\ on\ \mathbf{R}$ follows and this completes the proof of the theorem.

**COROLLARY.** Suppose $b_i a_{ij}$ (i, j = 1.2, … , n)in (1.1)are periodic with a common period say $T > 0$. *Then if* (1.2), (1.3)*and* (2.3)*hold the system* (1.1)*will have a periodic solution* $w(t) = \{w,(t), \ldots\ldots, w_n(t)\}$*of period* $T$ *such that*

$$w_1(t+T) = w_1(t), \qquad t \in \boldsymbol{R}: i = 1,2, \ldots, n.$$

**PROOF.** Proof follows from that of Theorem 2.1 if we Choose $T_k = KT(K = 0,1,2, \ldots)$*and note that*

$$x_1(t+KT) - w_1(t) \to 0\ as\ k \to \infty; i = 1,2, \ldots, n; t \in \boldsymbol{R},$$

$$x_i(t+(K+1)t) - w_1(t) \to 0 \ \ as \ \boldsymbol{K} \to \infty; i = 1,2,\dots,n; t \in \boldsymbol{R},$$

$$x_i(t+T+KT) - w_1(t+T) \to 0 \ \ as \ k \to \infty; i = 1,2,\dots,n; t \in \boldsymbol{R},$$

Implying that

$$w_1(t+T) = w_i(t); \quad i = 1,2,\dots,n; t \in \boldsymbol{R},$$

## 3. Global asymptotic stability

By definition we say that an almost-periodic solution say $w(t) = \{w_1(t),\dots,w_n(t)\} of (1.1) is\ globally\ asymptotically$ stable (or attractive) if and only if every other solution x(t)={ $x_1(t),\dots,x_n(t)$}with $x_1(o) > 0, i = 1.2,\dots,n$ is defined for all t□0 and satisfies

$$\lim_{t\to\infty} |w_1(t) - x_1(t)| = 0; \quad i = 1,2,\dots,n.$$

A consequence of such a global asymptotic stability is that there cannot be another strictly positive almost-periodic solution of (1.1)

**THEOREM 3.1** Suppose the almost- Periodic coefficients $b_1$, $a_{ij}$ (i, j = 1,2, … , n)

Satisfy (1.2)and (1.3)and the following

$$\inf_{t\in R} a_{ij}(t) > \left\{ \sum_{\substack{J=1 \\ J\neq 1}}^{n} \sup a_{ij}(t) \right\} + \mu; \qquad i = 1,2,\dots,n$$

For some positive number $\mu\ Then$ (1.1)has an almost-periodic solution (with strictly positive components) which is globally asymptotically stable.

We will omit the proof since the proof is identical with that of Theorem 3.1 of [3]

Under condition (3.2) the almost –periodic solution $w(t) = \{w_1(t), \ldots . w_n(t)\}$ of (1.1)enjoys a somewhat strong from of stability even if the coefficients of (1.1)are subjected to a class of perturbations We note (see Yoshizawa [7]page17)that if a function b:R→ $R$ is almost periodic then H(b),the "hull of b" is defined as follows:

H(b)= g|g:**R**→ $\boldsymbol{R}$ $and\ g(t)\lim_{k\to\infty} b(t+T_k) for$ some sequence

$T_k$ of real numbers the limit being uniform in $t \in R$.

By Theorem 2.2 of Yoshizawa [7, page 10], members of the hull of almost periodic functions are themselves almost periodic.

A system of the from

$$\frac{dy_1(t)}{dt} = y_1(t)\left\{B_i(t) - \sum_{J-1}^{n} A_{ij}\ (t)y_1(t)\right\},$$

*i=1,2,…,n; t>*$t_0 \in R$,

Where

$$B_i \in H(b_i); (A_{ij}), \qquad \text{i, j} = 1,2, \ldots, \text{n},$$

is known as a system from the hull of (1.1) As a consequence of the definition of the hull in (3.3)and (3.5)we shall have

$$\inf_{t\in R} B_i(t) = \text{b}_i^l; \quad \sup_{t\in R} B_i(t) = \text{b}_i^l;$$

$$\inf_{t\in R} A_{ij}(t) = \text{b}_i^l; \quad \sup_{t\in R} A_{ij}(t) = \text{a}_{ij}^u;$$

and hence the coefficients of(3.4)- (3.5)satisfy the same conditions as (1.2),(1.3)and (3.2)It will then follow from theorems(2.1)and (3.1)that solutions of (3.4)-(3.5)with $y_1(t_0) > 0, i = 1,2, \ldots, n\ will\ have$ the following property:

$$y_1(t) > 0 \ \ for\ t \geq t_0; i = 1,2,\dots,n,$$

$$\sum_{i=1}^{n} y_1(t) remains\ bounded\ for\ t \geq t_0$$

**DEFINITION** A Solution $x(t) = \{x_1(t), \dots, x_n(t)\}, t \in t_0 \in \boldsymbol{R}$,of (1.1)is said to be stable under perturbations from the hull of (1.1)if the following holds :let $y(t) = \{(y_1(t), \dots, y_n(t)\}$ be any solution of (3.4)-(3.5)such that $y_1(t_0) > 0$, i=1,2,…,n,and for some positive constants $\beta_1, \beta_2 \beta_1, \leq y_1(t) \leq \beta_2, t \geq t_0, \ \ i = 1,2,\dots,n$ Then given any $\varepsilon > 0$,there exists a constant $\delta > 0$ such that

$$\sum_{j=1}^{n} |x_1(t_0) - y_1(t_0)| < \delta$$

$$\sum_{i=1}^{n} \sup_{t \in R} |b_1(1) - B_1(t)| < \delta \quad = \sum_{i=1}^{n} |x_1(t) - y_i(t)| < \varepsilon$$

$$\sum_{i=1}^{n} \sum_{j=1}^{n} | \quad \sup_{t \in R} |a_{ij}\ (t) - A_{ij}(t)| < \delta$$

For $t \geq t_0$.

We can now formulate the following result which shows that all solution of (1.1)are stable under perturbations from the hull of (1.1)

**THEOREM 3.2**. Assume that the conditions of Theorem (3.1) hold. Then every strictly positive solution of (1.1) ( including its unique almost-periodic solution) is stable under perturbations from the hull of (1.1).

PROOF. Let $x(t) = \{x_1(t), \dots, x_n(t)\}$ and $y(t) = \{y_1(t), \dots, y_n(t)\}$ be any two strictly positive solutions of (1.1) and (3.4) – (3.5) respectively for $t \geq t_0 \in R$. Let $\beta_1, \beta_2$ be constants such that

$\beta_1 \leq x_i(t) \leq \beta_2$ ,
$\beta_1 \leq y_i(t) \leq \beta_2$ , $i = 1,2,\dots,n; t \geq t_0$ (3.7)

Consider now a function $V(t) = V(x(t)y(t))$ defined by

$$V(x(t)y(t)) = \sum_{i=1}^{n} |\log x_1(t) - \log, y_i(t)|, \quad t \geq t_0 \in \text{R}. \tag{3.8}$$

It is easy to see that

$$\frac{1}{\beta_2} \sum_{i=1}^{n} |x_i(t) - y_i(t)| \leq V(x(t)y(t))$$

$$\leq \frac{1}{\beta_1} \sum_{i=1}^{n} |x_i(t) - y_i(t)|, \; t \geq t_0 \in \text{R}. \tag{3.9}$$

Calculating the right derivative $D^+V$ of $V$ and using (1.1) and (3.4) we derive (after simplification) that

$$D^+V\left(x(t), y(t)\right) \leq -\mu \sum_{i=1}^{n} |x_i(t) - y_i(t)| + \sum_{i=1}^{n} |b_i(t) - B_i(t)|$$

$$+\beta_2 \sum_{j=1}^{n} \sum_{i=1}^{n} |a_{ij}(t) - A_{ij}(t)|; \quad t \geq t_0,$$

$$\leq -\mu\, \beta_1\, V\left(x(t), y(t)\right) + \sum_{i=1}^{n} |b_i(t) - B_i(t)|$$

$$+ \beta_2 \sum_{i=1}^{n} \sum_{j=1}^{n} |a_{ij}(t) - A_{ij}(t)|; \quad t \geq t_0, \tag{3.10}$$

An integration of (3.9) over $[t_0, t]$ with an application of (3.8) leads to

$$\frac{1}{\beta_2}\sum_{i=1}^{n}|x_i(t)-y_i(t)| \leq V(x(t),y(t)) \leq V(x(t_0),y(t_0)) + \left[\sum_{i=1}^{n}\sup_{t\in R}|b_i(t)-B_i(t)| + \beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}|a_{ij}(t)-A_{ij}(t)|\right]\frac{1}{\mu\beta_1},$$

$$t \geq t_0,$$

And hence

$$\sum_{i=1}^{n}|x_1(t)-y_1(t)| \leq \left(\frac{\beta_2}{\beta_1}\right)\sum_{i=1}^{n}|x_1(t_0)-y_1(t_0)| + \frac{\beta_2}{\mu\beta_2}\left[\sum_{i=1}^{n}\sup_{t\in R}|b_i(t)-B_i(t)|\right.$$

The construction reason is because is criticized in a number of usage guides: The reason for the long delays was because the costs greatly exceeded the original estimates. One objection to this construction is based on its redundancy: the word because (literally, by cause) contains within it the meaning of reason; thus saying the reason is because is like saying “The cause is by cause,” which would never be said. A second objection is based on the claim that because can introduce only adverbial clauses and that reason is requires completion by a noun clause. Critics would substitute that for because in the offending construction: The reason for the long delays in completing the project was that the costs.

$$\left.+\beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}|a_{ij}(t)-A_{ij}(t)|\right]. \qquad (3.11)$$

Now for any $\varepsilon > 0$ if one chooses a $\delta > 0$ such that

$$\delta < \min\left\{\frac{\varepsilon}{3}\frac{\beta_1}{\beta_2}, \frac{\varepsilon}{3}\frac{\beta_1}{\beta^2}\mu\right\} \qquad (3.12)$$

Then (3..6) will be satisfied and this completes the proof

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# . Jia Li's Prescientiial Dynamics And Contential Staticity Of Accentuation Attrition Models Applied To Cosmological And Astrophysical And Quantum Mechanical Engineering Problems

**Abstract**: **Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics ET AL ). We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems. Predator prey model literally form bastion, pillar, post, stylobate, sentinel towards the end of dismasting many of the asymptotic and solutional behaviour in finance, economic, and as we shall see in cosmology, quantum computation and other fields.**

## MODEL FORMULATION

Let $p_1(t, a_1)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_1}{dt} = -b_1 Q_1 + g_1(g_1)Q_1$$

Where $b$ is the natural death rate in the absence of credit types and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_1(0) = 0, \quad g_1(P_1) > 0, \; P_1 \; 0, \qquad \lim_{P_1 \to \infty} g_1(P_1) = (K_1)_1 \; \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_1 + (\rho_a)_1 + D_1(a_1, P_1, Q_1)\rho_1 + 0, \qquad 0 < a_1 \; \infty,$$

$$\rho_1(t, 0) = \int_0^{\infty} n_1(a_1, P_1, Q_1)\rho_1(t_1, a_1)da_1$$

Where $D_1$ is the mortality function and $n_1$ is the fertility function.

Put $_1$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_1}{dt}\right]_1 = x_1\left[\frac{dQ_2}{dt}\right]_1 - x_2\left[\frac{dQ_1}{dt}\right]_1$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_1 - x_4\left[\frac{dQ_2}{dt}\right]_1$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_1 - x_6\left[\frac{dQ_3}{dt}\right]_1$$

Where $x_1, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_1$ and $g_1(P_1)Q_1$

Now when you add the term $(b_1Q_1)$ with $g_1(P_1)(Q_1)$ we are infect see is that $(b_1Q_1)$ is credited incrementally part of $g_1(P_1)(Q_1)$ or instant nearly, when subtracted from $\frac{dQ_1}{dt}$, it means that $g_1(P_1)(Q_1)$ is slowly eaten by $\frac{dQ_1}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_1 = \mu_1(a_1) + (f_1)_1(P_1)_1 + (f_2)_1(Q)_1 \quad (2)$$

$$n_1 = \beta_1(a_1)h_1(Q_1), \quad (3)$$

Where $\mu_1$the age- dependent debit rate whit is $\mu_1(a_1)$ 0 $\forall a > 0$; $(f_1)_1$ is the density- dependent death rate with $f_1(0) = 0$, $(f_1^1)_1(P_1)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_1)_1(P_1) = (k_1)_1$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_1(0) = 0, (f_2^1)_1(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_1(Q_1) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^1(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_1 + (\rho_a)_1 = -\mu_1(a_1) + (f_1)_1P_1) + \big(f_2)_1(Q_1)\big)(\rho_1)_1 \quad (4)$$

$$\rho_1(t_1, 0) = h_1(Q_1)\int_0^\infty \beta_1(a_1)\rho_1(t, a_1)da, \quad (5)$$

$$\frac{dQ_1}{dt}$$

$$= -b_1Q_1 + g_1(P_1)Q_1 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_1$ and $Q_1$ of (4-6) satisfy

$$0 \leq \rho_1(t, a_1) \leq \rho_1^*(t, a_1),$$

$$0 \leq Q_1(t) \leq Q_1^*(t),$$

Where $\rho_1^*(t, a_1)$ and $Q_1^*(t)$, are solutions of

$$(\rho_t^*)_1 + (\rho_a^*)_1 = -(\mu_1(a_1) + (f_1)_1(P_1^*) + (f_2)_1 Q_1^*))\, \rho_1^*, \quad (7)$$

$$\rho_1^*(t, 0) = \int_0^\infty \beta_1(a_1)\, \rho_1^*(t, a_1) da_1 \quad (8)$$

$$\frac{dQ_1^*}{dt} = -b_1 Q_1^* + g_1(P_1^*) Q_1^*, \quad (9)$$

With $P_1^*(t) = \int_0^\infty \rho_1^*(t, a_1) da_1$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_1^*(t)}{dt} = P_1^*(t)\big(r_1 - (f_1)(P_1^*) - (f_2)_1(Q^*)\big), \quad (10)$$

$$\frac{dQ_1^*(t)}{dt} = Q_1^*(t)(-b_1 + g_1(P_1^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_1(a) e^{-\int_0^a \mu_1(\tau) dt} e^{-ra}\, da_1 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_1^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(t)dt}\, da_1 > 1, \qquad (H1)$$

Which guarantees $r_1 > 0$.

*Remark.* If $\beta_1(a) = [\beta_0]_1 a_1 e^{-\alpha_1 a_1}, \mu_1(a_1) \equiv \mu_1$ then the solution of (12) is

$$r_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_1]_1(P_1) = 0, [f_2](Q_1) = \lambda_1 Q_1\ h_1(Q_1) = 1, \qquad and\ \ g_1(P_1) = c_1 P_1$$

$$[\rho_1]_1 + [\rho_a]_1 = -(\mu_1 + \lambda_1 Q_1)\rho_1$$

$$\rho_1(t,0) = \int_0^\infty [\beta_0]_1 a_1 e^{-\alpha_1 a_1}\rho_1(t,a_1)\, da_1$$

$$\frac{dQ_1}{dt} = -b_1 Q_1 + c_1 P_1 Q_1$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential

equations

$$\frac{dP_1}{dt} = \omega_1 P_1 - \lambda_1 P_1 Q_1$$

$$\frac{dQ_1}{dt} = b_1 Q_1 - c_1 P_1 Q_1$$

Where $\omega_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$ THE EQUILIBRIUM $(\rho_0)_1(a_1), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_1(a_1), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_1 + (\rho_a)_1 = -\big(\mu_1(a_1) + (f_1)_1(P_1)\big)\rho,$$

$$\rho_1(t,0) = \int_0^\infty \beta_1(a_1)\rho_1(t,a_1)da_1 \qquad (13)$$

The solution of (13) has the form of

$$\rho_1(t,a_1) = Ce_1^{r(t-a_1)} - \int_0^a (\tau)d\tau - \int_0^1 (f_1)_1\big(P_1(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_1(t) = C_1 e^{r_1 t - \int_0^t (f_1)_1(P_1(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_1^{-1}(r)$ if $k_1 > r$.

This is equivalent to

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{(k_1)_1 a_1} da_1 < 1.$$

In fact, if we define

$$R_1(s_1) = \int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau)d\tau} e^{-s_1 a_1}\, da_1$$

$R(s)$ is a decreasing function. $R_1(r_1) = 1$ And $R_1(k_1)_1 < 1$ imply $r_1 < (k_1)_1$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_1(a_1), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{k_1 a_1} da_1 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_1 \equiv \int_0^\infty (\rho_0)_{11}(a_1)da_1$ *satisfies*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{-(f_1)_1(P_0)a_1} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_1(a_1), 0)$

Let $u_1(t, a_1) = \rho(t, a_1) - \rho_0(a_1)$ Then

$$N_1(t) \equiv \int_0^\infty u_1(t, a_1)da_1 = \int_0^\infty \rho_1(t, a_1)da_1 - \int_0^\infty (\rho_0)_1(a_1)da = P_1(t) - P_0)_1$$

And

$$(u_1)_1 + (u_a)_1 = (\rho_t)_1 + (\rho_a)_1 - (\rho_0^1)_1(a_1)$$

$$= -\mu_1(a_1) + (f_1)_1(P_1) + (f_2)_1(Q_1))(u_1 + (\rho_0)_1) + \mu_1(a_1)(\rho_0)_1 + (f_1)_1(P_0)_1(\rho_0)_1$$

$$= -\mu_1(a_1)u_1 - (f_1)_1(P_1) + (f_2)_1(Q_1))u_1 - (\rho_0)_1(f_1)_1(P_1) - (f_1)_1(P_0)_1$$

$$+ (f_2)_1(Q_1)), \qquad (15)$$

Where $(\rho_0^1)_1(a_1)$ denotes $d(\rho_0)_1(a_1)/da_1$

Expansion $(f_1)_1(P_1)$ and $(f_2)_1(Q_1)$ at $(P_0)_1$ and 0 yield respectively

$$(f_1)_1(P_1) = (f_1)_1((P_0)_1 + N_1)$$

$$= (f_1)_1(P_0)_1 + (f_1^1)_1(P_0)_1N_1 + o_1(N_1), \qquad (16)$$

$$(f_2)_1(Q_1)$$

$$= (f_2^1)_1(0)Q + o_1(Q_1) \qquad (17)$$

Write $f_1(P_0) = f_{10}, f_1'(P_0) = f_{10}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_1)_1 + (u_a)_1 \approx -(\mu_1(a_1) + (f_{10})_1)u - (f_{10}^1)_1(\rho_0)_1(N_1) - (f_{20}^1)_1(\rho_0)_1(Q_1)$$

$$u_1(t,0) = \rho_1(t,0)(\rho_0)_1(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t,a) + (\rho_0)_1(a_1)da_1 - (\rho_0)_1(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t,a_1)da_1 + (h_1(Q_1) - 1)(\rho_0)_1(0) \qquad (18)$$

Expanding

$$h_1(Q_1) = 1 + (h^1)_1(0)Q_1 + 0(Q_1) \equiv 1 + h_0'Q_1 + 0Q_1$$

And substituting it into (18), we arrive at

$$u_1(t,0) \approx \int_0^\infty \beta_1(a_1)u_1(t,a_1)da_1 + (h_{a_1}')(\rho_0)_1Q_1$$

The linearization is the following:

$$[u_1] + [u_1] = -(\mu_1(a_1) + [t_{10}]_1u_1 - [\alpha_1]_1(a_1)N_1 - [\alpha_2]_1(a_1)Q_1$$

$$u_1(t,0) = \int_0^\infty \beta_1(a_1)u_1(t,a_1)da_1 + [\xi_1]Q_1 \qquad (19)$$

$$\frac{dQ_1}{dt} = -b_1Q_1$$

where

$$[\alpha_1]_1(a_1) = [t_1^1 a]_1[p, o]_1(a_1)d2(a_1) = [f_{20}^1]_1(\rho\, o\,)_1(a_1)\xi = [\mathrm{h\ o}^1]_1[\rho \mathrm{o}]_1(\mathrm{o}) \text{ and } \mathrm{b}_1 = \mathrm{b}_1 - g_1(P_1)$$

If $g_1(P_0) > b$, clearly, $(\rho_0)_1(a_1), 0)$ is unstable and the predator population is persistent. Assume $g_1(P_0)_1 < b$.

Then, $Q_1(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_1 + u_a = -(\mu_1(a_1) + (f_{10})_1)u_1 - (\alpha_1)_1(a_1)N_1$$

$$= -(\mu_1(a_1) + f_{10})u_1 - \alpha_1(a_1)\int_0^\infty u_1(t, a_1)da_1$$

$$u_1(t, 0) = \int_0^\infty \beta_1\,(a_1)u_1(t, a_1)da_1$$

Using the method developed by Cushing [2, 3], we suppose

$$u_1(t, a_1) = \tilde{\mathrm{u}}_1(a_1)^{e^{\gamma(t-a_1)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_1}\, da_1$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_1)$ gives,

$$\tilde{\mathrm{u}}_1(a_1) = e^{-\int_0^a \mu_1(\tau_1)d\tau - (f_{10})_1 a}\left(\tilde{\mathrm{u}}_1(0) - w\int_0^a e^{\int_0^\tau \mu_1(s_1)ds + (f_{10})}\tau_1\,(\alpha_1)_1(\tau_1)e^{\gamma_1\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_1(a_1)e^{-\gamma_1 a_1 - (f_{10})_1 a_1 - \int_0^a \mu_1(\tau_1)d\tau}\right)\tilde{\mathrm{u}}_1(0)$$

$$+w\int_0^{\infty}\beta_1(a_1)e^{-\gamma_1 a_1-(f_{10})_1 a_1-\int_0^a \mu_1(\tau_1)d\tau}\int_0^a(\alpha_1)_1(\tau_1)e^{-\gamma_1\tau_1}e^{(f_{10})_1\tau_1+\int_0^{\tau}\mu_1(s_1)ds}d\tau_1\,da_1$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty}e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}da\right)u(0)$$

$$+\left(-1-\int_0^{\infty}e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}\int_0^a e^{f_{10}\tau+\int_0^{\tau}\mu(s)ds+y\tau}\alpha_1(\tau)d\tau\,da\right)w=0.\ (23)$$

Define $Y(\gamma,a)=e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta\,(a)Y(\gamma,a)da\right)u(0)$$

$$+w\int_0^{\infty}\beta(a)Y\,(\gamma,a)\int_0^a\alpha_1(\tau)Y^1(\gamma,\tau)d\tau\,da=0,(24)$$

$$\left(\int_0^{\infty}Y(\gamma,a)da\right)u(0)$$

$$+\left(-1-\int_0^{\infty}Y(\gamma,a)\int_0^a\alpha_1(\tau)Y^{-1}(\gamma,\tau)d\tau\,da\right)w=0.$$

Equation (20) has a solution of the form of $u(t,a)=u(a)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $u(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I(y,a)=\int_0^a\rho_0(\tau)Y^{-1}\,(\gamma,\tau)d\tau$. We obtain the characteristic equation

$$f_{10}^{l}\left(\int_0^\infty Y(\gamma,a)da\right)\left(\int_0^\infty \beta(a)Y(\gamma,a)I(\gamma,a)da\right)$$

$$+\left(1-\int_0^\infty \beta(a)Y(\gamma,a)da\right)\left(1+f_{10}^{l}\int_0^\infty Y(\gamma,a)I(\gamma,a)da\right)=0. \quad (25)$$

Since

$$\rho_0(a)=\rho_0(0)e^{-\int_0^a \mu(\tau)d\tau - f_{10}a}$$

$$I(\gamma,a)=\int_0^a \rho_0(\tau)e^{\gamma\tau+f_{10}\tau+\int_0^\tau \mu(s)ds}d\tau$$

$$=\frac{\rho_0(0)}{\gamma}(e^{\gamma a}-1)$$

Let

$$E(a)=e^{-f_{10}a-\int_0^a \mu(\tau)d\tau}$$

Then $[Y_1]_1[\gamma_1]_1,a_1)=\ E_1(a_1)e^{-\gamma_1 a_1}$. by a simple calculation,

$$\int_0^\infty \beta_1(a_1)Y_1(\gamma_1,a_1)I(\gamma_1,a_1)da=\frac{[\rho_0]_1(0)}{\gamma}\left(1-\int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}\ da\right);$$

$$\int_0^\infty Y_1(\gamma_1,a_1)I(\gamma_1,a_1)da=\frac{[\rho_0]_1(0)}{\gamma_1}\left(\int_0^\infty E_1(a_1)e^{-\gamma_1 a_1}\ da_1\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_1(0)}{\gamma_1}[f_{10}^{1}]\int_0^\infty E_1(a_1)e^{-\gamma_1 a_1}da_1\ (\int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1)$$

$$+\left(1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right).$$

$$\left(1+\frac{(\rho_0)_1(0)}{\gamma_1}(f_{10}^{'})_1\left(\int_0^{\infty}E_1(a_1)da_1-\int_0^{\infty}E_1(a_1)e^{-\gamma_1 a_1}da_1\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right)(1+\frac{(f_{10}^{'})_1}{\gamma_1}(\rho_0)_1(0)\int_0^{\infty}E_1(a_1)da_1)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_1(a_1)Y_1(\gamma_1 a_1)da_1=1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\neq 0,$$

Since $\int_0^{\infty}\beta_1(a_1)Y_1(\gamma_1 a_1)\int_0^{a}\alpha_1(\tau_1)Y^{-1}(\gamma_1,\tau_1)d\tau_1 da_1>0$. Therefore, the characteristic solution is

$$\gamma_1=-(f_{10}^{'})_1(\rho_0)_1\int_0^{\infty}E_1(a_1)da_1.$$

Since $f_1'(P_0)_1>0$, we obtain the following result

THEOREM 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_1),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_1),0)$ which is related to the extinction of debit summation as has been referred to in the foregoing and to The possible persistence of the prey

population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_1 + (\mu_1(a_1) + [f_1]_1(P_1^*) + [f_2]_1(Q_1^*))P_1^* = 0, \quad (26)$$

$$p_1^*(0) = h_1(Q_1^*)\int_0^\infty \beta_1(a_1)p_1^*(a_J)p_1^*(a_1)da_1 \quad (27)$$

$$g_1(P_1^*) - b_1 = 0, \quad (28)$$

Where $P_1^* = \int_0^\infty p_1^*(a_1)\, da_1$

Equation (28) yields $P^* = g^{-1}(b)$,

From (26),

$$p_1^*(a) = p_1^*(0)e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}$$

Substituting into (27), we need to solve

$$\left(1 - h_1(Q_1^*)\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}da_1\right)p_1^*(0) = 0$$

For $Q_1^*$

Fix $p_1^*(0)$ and $P_1^*$ Let

$$H_1(z_1) = 1 - h_1(z_1)\int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1}da_{,1}$$

And assume that

$$[k_3]_1 \int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(g_1^{-1}(b_1))a_1 - [k_2]_1 a_1}\, da_1 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty} [H_1(z_1)] > 0.$$

Clearly,

$$\frac{dH_1}{dz_1} = h_1(z)(\,[f_2']_1(z_1)\,) \int_0^{\infty} a_1\beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - f_1(P_1^*)a_1 - f_2(z_1)a_1} da_1$$

$-h'(z_1) \int_0^{\infty} \beta_1(a_1)\, e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1}\ da_1$

By the assumptions on $\left[f_2^1\right]_1 and\ [h^1]^1 dH_1/dz > 0$. Hence, if

$$H_1(0) = 1 - \int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1} da_1 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_1(0) < 0$ follows if $g_1^-(b_1) < [f_1^{-1}]_1(r_1)$, where $r_1$ is the solution of (12).

Indeed, since

$$R_1(s_1) = \int_0^{\infty} \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 e^{-y_1^a}} da_1$$

Is a decreasing function of $s_1$ when $P_1^* = [g^{-1}]_1(b_1) < [f_1^{-1}]_1(r_1)$, that is $[f_1]_1(P_1^*) < r$, $R\big(f_1(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_1(g^{-1}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1\left(g_1^{-1}(b)\right)a_1} > 1. \qquad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_1^*(a_1), Q_1^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_1^*(a_1), Q_1^*)$ is more complicated. We again consider perturbations $m_1(t_1 a_1) = p_1(t_1, a_1) - p_1^*(a_1)$ and $W_1(t) = Q_1(t) - Q^*$.Then

$$M_1(t) \equiv \int_0^\infty m_1(t, a_1)da_1 = P_1(t) - P_1^*,$$

$$[m_1]_1 + [m_a]_1 = [p_1]_1 + [p_a]_1 - \left[p^*\right]_1(a_1) \approx -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1$$

$$-[f_{1*}^1]_1 p_1^*(a_1)M - [f_{2*}^1]_1 p_1^*(\mathrm{a})W_1,$$

$$m_1(t,0)p_1(t,0) - p_1^*(0) \approx h_* \int_0^\infty \beta_1(a_1)m_1(t,a_1)\,da + \frac{[h_*^1]_1}{[h_*]_1}\, p_1^*(0)W_1$$

$$\frac{dW_1}{dt} = \frac{dQ_1}{dt} \approx [g_*^1]_1 Q_1^* M_{,1}$$

Where $[f_{1*}]_1$ denoting $[f_1(P^*)]_1$, $[f_{1*}]_1$ denoting $[f_2(P^*)]_1$, $\left[f_{1*}^{'}\right]_1$ denoting $\left[f^{'}{}_1(P^*)\right]_1 \left[f_{2*}^{'}\right]_1$ denoting

$[f_2]_1\,(Q_1^*,\quad h_1^*$ denoting $h_*^{'}$ denoting $h_1^{'}\,(Q_1^*)$, and $\left[g_*^{'}\right]_1$

Denoting $\left[g^{'}\right]_1(P_1^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_1]_1 + [m_a]_1 = -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1 - [f'_{1*}]_1[p_1^*](a_1)M_1 - [f'_{2*}]_1 p_1^*(a_1)W_1$$

$$m_1(t,0) = [h_*]_1 \int_0^\infty \beta_1(a_1)m_1(t,a_1)da_1 + \frac{[h_*^1]_1}{[h^*]_1}[p^*]_1(0)\,W_{,1} \qquad (30)$$

$$\frac{dW_1}{dt} = [g_*^1]_1[Q_1^*]M_1$$

Integrating $[m_1]_1(t,a_1)$ with respect to a from 0 to $\infty$, we have

$$\frac{dM_1(t)}{dt} = \int_0^\infty [mt]_1(t,a_1)da_1$$

$$= -\int_0^\infty [m_a]_1\,(t,a_1)da_1 - \int_0^\infty \mu(a_1)m(t,a_1)da_1$$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f'_{1*}]_1 \int_0^\infty p_1^*(a_1)da_1\right)M_1 - [f_{2*}^1]_1 \int_0^\infty p_1^*(a_1)da_1 W_1$$

$$= m_1(t,0) - \int_0^\infty \mu_1(a_1)\,m_1(t,a_1)da_1$$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f_{1*}']_1 P_1^*\right)M_1 - [f_{2*}']_1 P_1^* W_{,1}$$

From (30), we arrive at the equations

$$\frac{dW_1(t)}{dt} = \int_0^\infty ([h_*]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1$$

$$-\left([f_{1*}] + [f_{2*}]_1 + [f_{1*}']_1 P_1^*\right)M_1 + \left(\frac{[h_*']_1}{[h_*]^1}\,P_1^*(0) - f_{2*}' P_1^*\right)W_1,$$

(31)

$$\frac{dW_1(t)}{dt}\left[g^{'}{}_{*}\right]_1 Q_1^{*} M_{.1}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{*}]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1 = ([h_{*}]_1\beta_1(\theta_1) - \mu_1(\theta_1))M_{,1}$$

Where $0< \theta_1 = \theta_1(t) < \infty$.

And if $a_1^2\,[a_1^2]_1<4\,[a_2]_1<4\,a_2 e$,

Re $\lambda_1\left[\lambda_1\right]_1 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_1 = -\frac{1}{2}a_1[a2\ ]_1 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_{*}\,[h_{*}]_1\beta_1(a_1) - \mu_1(a_1)\right|$. Then$|n_1(t)|$ ⊠Δ∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium($\rho^{*}(a), \boldsymbol{Q}^{*}$) is asymptotically stable if either $a_1^2 <4\,a_2 e$ and $[a_1] > 2\Delta$,or $[a_1^2]$ ⊠4$[a_2]$e and $(a_1 - \sqrt{(a_1^2)_1 - 4(a_2)_1 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+1}\right)$ and $\beta_1(a_1) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_1(a_1)\mathrm{e}^{-\int_0^{\mathrm{a}} \mu_1(\tau_1)+d\tau_1 \mathrm{da}} = \frac{1}{(\mathrm{m}+\mu_0+1)^2} + \frac{2}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_1 + [\mu_0]$ 0.5. Assume$[f_1](x_1)= K_1 x_1$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_1(x_1)= [0x_1]/(1 + x_1)$; the functional responses $[[f_2](x) = [0]\, x_1/(1 + x_1)$and $h_1(x_1)=(1+\omega_1 x_1)/(1 + x_1)$(See Holling [7] and Rashevsky [13])

Equation (14),now has the form of

$$\frac{m_1 + [\mu_0] + [\mathrm{K}][P_0])}{[m_1]_1 + (\mu_0)_1 + [\mathrm{K}][P_0])_1^3\ [\mathrm{K}]_1[P_0]_1^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_1$i e $[P_0] < b/(\delta_1 - b_1)$then the equilibrium$(P_0)\ [P_0]_1(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_1$

$[f_{1*}] = K_1\ \mathrm{P}^*\ [f_{2*}] = \frac{\sigma \mathrm{Q}_1^*}{1+\mathrm{Q}_1^*} \quad [h_*] = \frac{1+\omega_1 \mathrm{Q}_1^*}{1+\mathrm{Q}_1^*}$

$[f_{1*}] = K_1\ [\mathrm{f}_{2*}^1] = \frac{\sigma_1}{(1+\mathrm{Q}_1^*)^2}[\ \mathrm{g}_*^1] = \frac{\sigma_1}{(1+\mathrm{P}_1^*)^2} \quad [\mathrm{h}_*^1] = \frac{\omega_1 - 1}{(1+\mathrm{Q}_1^*)^2}$

## AN ANALOGY BETWEEN PREDATOR PREY AND ACCENTUATION ATTRION MODELS

**Best example how the model could be applied is the Schrödinger's equation itself wherein RHS gobbles up LHS, and the terms on RHS do the same with each other. It is like adding milk to water. Whether it happens instantaneously or time lag is there is a different question altogether dealt at the opportune points**

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| **Best example that can be given is the formulation of the quantum wave equation albeit in a nascent form for the recently found gravitational waves. This has been studied in length in various papers of gravitational waves, wherein the attribution of how the model is used is described. Note the calculation holds for only one module. This is in consideration to the fact that the relation is beteen the both debit systems and credit systems which as we have enucleated in the foregoing bears ample testimony for the application of the model** |
| **Gravitational quantum wave equation:** |
| **Equation Time-dependent equation (Wikipedia)** |
| **(1) In quantum mechanics, the Schrödinger equation is (=) a partial differential equation that describes how the quantum state of a quantum system changes with (e&eb) time. It was formulated in late 1925, and published in 1926, by the Austrian physicist Erwin Schrödinger.[1]** |
| **(2) In classical mechanics Newton's second law, (F = ma), is used to make (eb) a mathematical prediction as to what path a given system will take following a set of known initial conditions.** |
| **(3) In quantum mechanics, the analogue of Newton's law is (=) Schrödinger's equation for (e) a quantum system (usually atoms, molecules, and subatomic particles whether free, bound, or localised).** |
| **(4) It is not a simple algebraic equation, but in general a linear partial differential equation, describing (eb) the time-evolution of the system's wave function (also called a "state function").[2]:1–2** |
| **(5) The concept of a wavefunction is (=) a fundamental postulate of quantum mechanics.** |
| **(6) Although Schrödinger's equation is often presented as a separate postulate, some authors [3]: Chapter 3 show that some properties resulting from Schrödinger's equation may be deduced just from (e) symmetry principles alone, for example the commutation relations.** |
| **(7) Generally, "derivations" of the Schrödinger equation demonstrate (eb) its mathematical plausibility for describing (eb) wave–particle duality, but to date there are no universally accepted derivations of (e) Schrödinger's equation from (e) appropriate axioms.** |
| **(8) In the Copenhagen interpretation of quantum mechanics, the wave function is (=) the most complete description that can be given of (e) a physical system.** |
| **(9) Solutions to Schrödinger's equation describe (eb) not only molecular, atomic, and subatomic systems, but also macroscopic systems, possibly even the whole universe.[4]:292ff** |

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| **(10) The Schrödinger equation is consistent with (e&eb) classical mechanics but not with (e&eb) special relativity.** |
| **(11) Making quantum mechanics consistent with (e&eb) special relativity is one of the great achievements of (e) Quantum Field Theory and requires (e) the introduction of the creation and annihilation operators so that (e) the number of particles can no longer be considered constant.** |
| **(12) The Schrödinger equation is not the only way to make predictions in (eb) quantum mechanics—other formulations can be used, such as Werner Heisenberg's matrix mechanics, and Richard Feynman's path integral formulation** |
| **(13) The form of the Schrödinger equation depends on (=) the physical situation (see below for special cases).** |
| **(14) The most general form is the time-dependent Schrödinger equation, which gives a description of a system evolving with time: [5]:143** |
| **(15) A wave function that satisfies the non-relativistic Schrödinger equation with V = 0. In other words, this corresponds to (e &eb) a particle traveling freely through (e&eb) empty space. Time-dependent** |
| **Schrödinger equation (general)** |
| **(16) i ℏ ∂ ∂ t Ψ ( r , t ) = H ^ Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t) where i is the imaginary unit, ħ is the Planck constant divided by 2π, the symbol ∂/∂t indicates a partial derivative with respect to time t, Ψ (the Greek letter psi) is the wave function of (e) the quantum system, and Ĥ is the Hamiltonian operator (which characterizes the total energy of any given wave function and takes different forms depending on the situation).** |
| **(17) Each of these three rows is a wave function which satisfies (eb) the time-dependent Schrödinger equation for a harmonic oscillator. Left: The real part (blue) and imaginary part (red) of the wave function. Right:** |
| **(18) The probability distribution of finding the particle with (e&eb) this wave function at a given position.** |
| **(19) The top two rows are examples of stationary states, which correspond to (e&eb) standing waves.** |
| **(20) The bottom row is an example of a state which is (=) not a stationary state.** |
| **(21) The right column illustrates (eb) why stationary states are (=) called "stationary".** |

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| **(22) The most famous example is the non-relativistic Schrödinger equation for (e) a single particle moving in an electric field (but not a magnetic field; see the Pauli equation): [6]** |
| **Time-dependent Schrödinger equation** |
| **(Single non-relativistic particle)** |
| **(23) i ℏ ∂ ∂ t Ψ ( r , t ) = [ − ℏ 2 2 μ ∇ 2 + V ( r , t ) ] Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)** |
| **(24) Here μ is (=) the particle's "reduced mass", V is its potential energy, ∇2 is the Laplacian (a differential operator), and Ψ is (=) the wave function (more precisely, in this context, it is called the (=) "position-space wave function").** |
| **(25) In plain language, it means "total energy equals (=) kinetic energy plus potential energy", but the terms take unfamiliar forms for (e) reasons explained below.** |
| **(26) Given the particular differential operators involved, this is (=) a linear partial differential equation.** |
| **(27) It is also a diffusion equation, but unlike the heat equation, this one is (=) also a wave equation given (e&eb) the imaginary unit present in the transient term.** |
| **(28) The term "Schrödinger equation" can refer to (e) both the general equation (first box above), or the specific nonrelativistic version (second box above and variations thereof).** |
| **(29) The general equation is indeed quite general, used throughout quantum mechanics, for everything from the Dirac equation to quantum field theory, by plugging in various complicated expressions for the Hamiltonian. The specific nonrelativistic version is a simplified approximation to reality, which is quite accurate in many situations, but very inaccurate in others (see relativistic quantum mechanics and relativistic quantum field theory)** |
| **(30) To apply the Schrödinger equation, the Hamiltonian operator is set up for the system, accounting for the kinetic and potential energy of the particles constituting the system, then inserted into the Schrödinger equation. The resulting partial differential equation is solved for the wave function, which contains information about the system.** |
| **Time-independent equation** |

**(31) The time-dependent Schrödinger equation described above predicts that wave functions can form standing waves, called stationary states (also called "orbitals", as in atomic orbitals or molecular orbitals). These states are important in their own right, and if the stationary states are classified and understood, then it becomes easier to solve the time-dependent Schrödinger equation for any state. Stationary states can also be described by a simpler form of the Schrödinger equation, the time-independent Schrödinger equation. (This is only used when the Hamiltonian itself is not dependent on time explicitly. However, even in this case the total wave function still has a time dependency.)**

**Time-independent Schrödinger equation (general)**

**(32) E Ψ = H ^ ⊠ Ψ {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi } {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi }**

**In words, the equation states:**

**(33) When the Hamiltonian operator acts on a certain wave function Ψ, and the result is proportional to the same wave function Ψ, then Ψ is a stationary state, and the proportionality constant, E, is the energy of the state Ψ.**

**(34) The time-independent Schrödinger equation is discussed further below. In linear algebra terminology, this equation is an eigenvalue equation.**

**(35) As before, the most famous manifestation is the non-relativistic Schrödinger equation for a single particle moving in an electric field (but not a magnetic field):**

**Time-independent Schrödinger equation (single non-relativistic particle)**

**(36) E Ψ ( r ) = [ − ℏ 2 2 μ ∇ 2 + V ( r ) ] Ψ ( r ) {\displaystyle E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )} E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )**

**Gravitational waves: (Wikipedia)**

**(37) In theory, the loss of energy through gravitational radiation could eventually drop the Earth into the Sun. However, the total energy of the Earth orbiting the Sun (kinetic energy + gravitational potential energy) is about 1.14×1036 joules of which only 200 Watts (joules per second) is lost through gravitational radiation, leading to a decay in the orbit by about 1×10−15 meters per day or roughly the diameter of a proton. At this rate, it would take the Earth approximately 1×1013 times more than the current age of the Universe to spiral onto the Sun. This estimate overlooks the decrease in r over time, but the majority of the time the bodies are far apart and only radiating slowly, so the difference is unimportant in this example.**

**More generally, the rate of orbital decay can be approximated by[39]**

**(38) d r d t = − 64 5 G 3 c 5 ( m 1 m 2 ) ( m 1 + m 2 ) r 3 , {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,} {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,}**

**where r is the separation between the bodies, t time, G Newton's constant, c the speed of light, and m1 and m2 the masses of the bodies. This leads to an expected time to merger of [39]**

**t = 5 256 c 5 G 3 r 4 ( m 1 m 2 ) ( m 1 + m 2 ) . {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.} {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.}**

**(39) For example, a pair of solar mass neutron stars in a circular orbit at a separation of 1.89×108 m (189,000 km) has an orbital period of 1,000 seconds, and an expected lifetime of 1.30×1013 seconds or about 414,000 years. Such a system could be observed by LISA if it were not too far away. A far greater number of white dwarf binaries exist with orbital periods in this range. White dwarf binaries have masses in the order of the Sun, and diameters in the order of the Earth. They cannot get much closer together than 10,000 km before they will merge and explode in a supernova which would also end the emission of gravitational waves. Until then, their gravitational radiation would be comparable to that of a neutron star binary.**

**(40) By substituting value of dr/dt in Schrödinger's wave equation we get the quantum wave equation for gravitational waves.**

In order to have (H3) and (H4) hold, $m, \mu_0$, K, b, δ, and $\omega$ need to satisfy

$$\frac{\omega_1}{(m+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1+b_1}+\sigma)^2}+\frac{2\omega_1}{([m_1]_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1+b_1}+\sigma_1)^3}<1$$

And

$$\frac{1}{(m_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1-b_1})^2}+\frac{2\omega_1}{(m_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1-b_1})^3}>1$$

When the inequalities (33) and (34) hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$(m_1 + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 - b_1} + (m_1 + [\mu_0]_1 + \frac{K_1 b_1}{\delta_1 - b_1} + \sigma_1)\, Q_1^*)^3$

$= m_1 + [\mu_0]_1 + 3 + \frac{K_1 b_1}{\delta_1 - b_1}(1 + Q^*)(1 + \omega_1 Q_1^*$

Now

$$\left[a_1^2\right] - 4\left[a_2\right] e = \left(\frac{2K_1 b_1}{\delta_1 - b_1}\right)^2 + \frac{4K_1 b_1 \sigma_1 Q_1^*}{(\delta_1 - b_1)(1 + Q_1^*)} + \frac{\sigma_1^2 Q_1^{*2}}{(1 + Q_1^*)^2}$$

$$\frac{4\, b_1 \sigma_1 (\delta_1 - b) Q^*}{\delta_1 (1 + Q_1^*)^2} - \frac{4\, (\delta_1 - b_1)^2 (1 - \omega_1) Q_1^*}{\delta_1 (1 + \omega_1 Q_1^*)(1 + Q_1^*)} P^*(0)$$

And

$\Delta = \lim_{0 < a < \infty} |[h_*]_1\, a e_1^{mu}\ [\mu_0]_1 - \frac{a_1}{a_1 + 1}$

Hence if $[a_1^2 > 4(a_2)_1 e$ and

$$[a_1] = \frac{2K_1 b_1}{\delta_1 - b_1} + \frac{\sigma_1 Q_1^*}{1 + Q_1^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_1]_1 = [f_{1*}]_1 + [f_{2*}]_1 + [f_{1*}] P_1^*$$

$$[a_2]_1 = \left[f^{1}_{2*}\right] P^{*}_{1} - \frac{h^{1}_{*}}{h_{*}} P^{*}(0)$$

$$e=[g^{1}_{*}]_1 Q^{*}_{1}$$

$$\eta_1(t) = \left[h_{*}\right]_1 \beta_1 \left(\theta_1(t)\right) - \mu_1(\theta(t)),$$

The system(31)can be written as the following system:

$$[M(t)]\left[a_1\right]M_1 - [a_2]_1 W_1 - [n(t)]M_1$$

$$w_1(t) = eM_1$$

Write

$$A=\begin{pmatrix} -[a_1]_1 - [a_2]_1 \\ e \qquad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_1 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_1 \\ W_1 \end{pmatrix}$$

The solution of (32)is

$$X_1(t) = e^{A,t}[X]_1(0) + \int_0 e^{At}$$

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# Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

## (Module Three)

**Abstract: Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-435) 1990. We demonstrate its applicability cosmology/ quantum computation economics et al. As stated in the foregoing we shall apply model to various cosmological problem and show the relevance to the same. . To exemplify model can be applied to the following situational significations: Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an**

**outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 Masaru Shibata and Takashi Nakamura** perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. **Authors** use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society **Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995**

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(g_3)Q_3$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_3(0) = 0, \quad g_3(P_3) > 0, \quad P_3 \text{ 0}, \qquad \lim_{P_3 \to \infty} g_3(P_3) = (K_3)_3 \text{ ∞}. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_3 + (\rho_a)_3 + D_3(a_3, P_3, Q_3)\rho_3 + 0, \qquad 0 < a_3 \text{ ∞},$$

$$\rho_3(t,0) = \int_0^\infty n_3(a_3, P_3, Q_3)\rho_3(t_3, a_3)da_3$$

Where $D_3$ is the mortality function and $n_3$ is the fertility function.

Put $_3$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_3}{dt}\right]_3 = x_3\left[\frac{dQ_2}{dt}\right]_3 - x_2\left[\frac{dQ_3}{dt}\right]_3$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_3 - x_4\left[\frac{dQ_2}{dt}\right]_3$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_3 - x_6\left[\frac{dQ_3}{dt}\right]_3$$

Where $x_3, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_3$ and $g_3(P_3)Q_3$

Now when you add the term $(b_3Q_3)$ with $g_3(P_3)(Q_3)$ we are infect see is that $(b_3Q_3)$ is credited incrementally part of $g_3(P_3)(Q_3)$ or instant nearly, when subtracted from $\frac{dQ_3}{dt}$ , it means that $g_3(P_3)(Q_3)$ is slowly eaten by $\frac{dQ_3}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_3 = \mu_3(a_3) + (f_3)_3(P_3)_3 + (f_2)_3(Q)_3 \quad (2)$$

$$n_3 = \beta_3(a_3)h_3(Q_3), \quad (3)$$

Where $\mu_3$the age- dependent debit rate whit is $\mu_3(a_3)$ 0 $\forall a > 0$; $(f_3)_3$ is the density- dependent death rate with $f_3(0) = 0$, $(f_3^3)_3(P_3)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_3)_3(P_3) = (k_3)_3$ ; $f_2$ is the functional response of

creditors to prey with $[f_2]_3(0) = 0, (f_2^3)_3(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_3\ (Q_3) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 3, h^3(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_3 + (\rho_a)_3 = -\mu_3(a_3) + (f_3)_3 P_3) + \big(f_2)_3(Q_3)\big)(\rho_3)_3 \quad (4)$$

$$\rho_3(t_3, 0) = h_3(Q_3)\int_0^\infty \beta_3\,(a_3)\rho_3(t, a_3)da, \quad (5)$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(P_3)Q_3 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_3$ and $Q_3$ of (4-6) satisfy

$$0 \ \rho_3(t, a_3) \ \rho_3^*(t, a_3),$$

$$0 \ Q_3(t) \ Q_3^*(t),$$

Where $\rho_3^*(t, a_3)$ and $Q_3^*(t)$, are solutions of

$$(\rho_t^*)_3 + (\rho_a^*)_3 = -(\mu_3(a_3) + (f_3)_3(P_3^*) + (f_2)_3 Q_3^*))\,\rho_3^*, \quad (7)$$

$$\rho_3^*\,(t, 0) = \int_0^\infty \beta_3(a_3)\,\rho_3^*\,(t, a_3)da_3 \quad (8)$$

$$\frac{dQ_3^*}{dt} = -\,b_3\,Q_3^* + g_3(P_3^*)Q_3^*, \quad (9)$$

With $P_3^*(t) = \int_0^\infty \rho_3^*\ (t, a_3)da_3$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_3^*(t)}{dt} = P_3^*(t)\big(r_3 - (f_3)(P_3^*) - (f_2)_3(Q^*)\big), \quad (10)$$

$$\frac{dQ_3^*(t)}{dt} = Q_3^*(t)(-b_3 + g_3(P_3^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_3(a)e^{-\int_0^a \mu_3(\tau)dt}e^{-ra}\ da_3 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_3^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(t)dt}\ da_3 > 1, \quad (H1)$$

Which guarantees $r_3 > 0$.

*Remark.* If $\beta_3(a) = [\beta_0]_3 a_3 e^{-\alpha_3 a_3}, \mu_3(a_3) \equiv \mu_3$ then the solution of (12) is

$$r_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_3]_3(P_3) = 0, [f_2](Q_3) = \lambda_3 Q_3\ h_3(Q_3) = 1, \quad and\ \ g_3(P_3) = c_3 P_3$$

$$[\rho_3]_3 + [\rho_a]_3 = -(\mu_3 + \lambda_3 Q_3)\rho_3$$

$$\rho_3(t,0) = \int_0^{\infty} [\beta_0]_3 a_3 e^{-\alpha_3 a_3} \rho_3(t,a_3)\, da_3$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + c_3 P_3 Q_3$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_3}{dt} = \omega_3 P_3 - \lambda_3 P_3 Q_3$$

$$\frac{dQ_3}{dt} = b_3 Q_3 - c_3 P_3 Q_3$$

Where $\omega_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$

## THE EQUILIBRIUM $(\rho_0)_3(a_3), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_3(a_3), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_3 + (\rho_a)_3 = -\big(\mu_3(a_3) + (f_3)_3(P_3)\big)\rho,$$

$$\rho_3(t,0)$$

$$= \int_0^\infty \beta_3\,(a_3)\rho_3(t,a_3)da_3 \qquad (13)$$

The solution of (13) has the form of

$$\rho_3(t,a_3) = Ce_3^{r(t-a_3)} - \int_0^a (\tau)d\tau - \int_0^3 (f_3)_3\big(P_3(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_3(t) = C_3 e^{r_3 t - \int_0^t (f_3)_3(P_3(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_3^{-3}(r)$ if $k_3 > r$. This is equivalent to

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{(k_3)_3 a_3} da_3 < 1.$$

In fact, if we define

$$R_3(s_3) = \int_0^\infty \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau)d\tau} e^{-s_3 a_3}\, da_3$$

$R(s)$ is a decreasing function. $R_3(r_3) = 1$ And $R_3(k_3)_3 < 1$ imply $r_3 < (k_3)_3$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_3(a_3), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{k_3 a_3} da_3$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_3 \equiv \int_0^\infty (\rho_0)_{33}(a_3)da_3$ *satisfies*

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{-(f_3)_3(P_0)a_3} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_3(a_3), 0)$

Let $u_3(t, a_3) = \rho(t, a_3) - \rho_0(a_3)$ Then

$$N_3(t) \equiv \int_0^{\infty} u_3(t, a_3) da_3 = \int_0^{\infty} \rho_3(t, a_3) da_3 - \int_0^{\infty} (\rho_0)_3(a_3) da = P_3(t) - P_0)_3$$

And

$$(u_3)_3 + (u_a)_3 = (\rho_t)_3 + (\rho_a)_3 - (\rho_0^3)_3(a_3)$$

$$= -\mu_3(a_3) + (f_3)_3(P_3) + (f_2)_3(Q_3))(u_3 + (\rho_0)_3) + \mu_3(a_3)(\rho_0)_3 + (f_3)_3(P_0)_3(\rho_0)_3$$

$$= -\mu_3(a_3)u_3 - (f_3)_3(P_3) + (f_2)_3(Q_3))u_3 - (\rho_0)_3 (f_3)_3(P_3) - (f_3)_3(P_0)_3$$

$$+ (f_2)_3(Q_3)), \qquad (15)$$

Where $(\rho_0^3)_3(a_3)$ denotes $d(\rho_0)_3(a_3)/da_3$

Expansion $(f_3)_3(P_3)$ and $(f_2)_3(Q_3)$ at $(P_0)_3$ and 0 yield respectively

$$(f_3)_3(P_3) = (f_3)_3((P_0)_3 + N_3)$$

$$= (f_3)_3(P_0)_3 + (f_3^3)_3(P_0)_3 N_3 + o_3(N_3), \qquad (16)$$

$$(f_2)_3(Q_3)$$

$$= (f_2^3)_3(0)Q + o_3(Q_3) \qquad (17)$$

Write $f_3(P_0) = f_{30}, f_3'(P_0) = f_{30}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_3)_3 + (u_a)_3 \approx -(\mu_3(a_3) + (f_{30})_3)u - (f_{30}^3)_3(\rho_0)_3(N_3) - (f_{20}^3)_3(\rho_0)_3(Q_3)$$

$$u_3(t, 0) = \rho_3(t, 0)(\rho_0)_3(0)$$

$$= h_3(Q_3)\int_0^{\infty} \beta_3(a_3)(u_3(t,a) + (\rho_0)_3(a_3)da_3 - (\rho_0)_3\,(0)$$

$$= h_3(Q_3)\int_0^{\infty} \beta_3(a_3)(u_3(t,a_3)da_3 + (h_3(Q_3) - 3)(\rho_0)_3\,(0) \qquad (18)$$

Expanding

$$h_3(Q_3) = 1 + \ (h^3)_3(0)Q_3 + 0(Q_3) \equiv 1 + h_0^{'}Q_3 + 0Q_3$$

And substituting it into (38), we arrive at

$$u_3(t,0) \approx \int_0^{\infty} \beta_3\ (a_3)u_3(t,a_3)da_3 + (h_{a_3}^{'})(\rho_0)_3Q_3$$

The linearization is the following :

$$[u_3] + [u_3] = -(\mu_3(a_3) + [t_{30}]_3u_3 - [\alpha_3]_3(a_3)N_3 - [\alpha_2]_3(a_3)Q_3$$

$$u_3(t,0) = \int_0^{\infty} \beta_3(a_3)u_3(t,a_3)da_3 + [\xi_3]Q_3 \qquad (19)$$

$$\frac{dQ_3}{dt} = -b_3Q_3$$

where

$$[\alpha_3]_3(a_3) = [t_3^3a]_3[p,o]_3(a_3)d2(a_3) = [f_{20}^3]_3(\rho\ o\ )_3(a_3)\xi = [\mathrm{h\ o}^3]_3[\rho\mathrm{o}]_3(\mathrm{o}) \text{ and } \mathrm{b}_3 = \mathrm{b}_3 - g_3(P_3)$$

If $g_3(P_0) > b$, clearly, $(\rho_0)_3(a_3), 0)$ is unstable and the predator population is persistent. Assume $g_3(P_0)_3 < b$.

Then, $Q_3(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_3 + u_a = -(\mu_3(a_3) + (f_{30})_3)u_3 - (\alpha_3)_3(a_3)N_3$$

$$= -(\mu_3(a_3) + f_{30})u_3 - \alpha_3(a_3)\int_0^\infty u_3(t, a_3)da_3$$

$$u_3(t, 0) = \int_0^\infty \beta_3\,(a_3)u_3(t, a_3)da_3$$

Using the method developed by Cushing [2, 3], we suppose

$$u_3(t, a_3) = \tilde{u}_3(a_3)^{e^{\gamma(t-a_3)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_3}\,da_3$. Then,

A solution of the first differential equation for $\tilde{u}(a_3)$ gives,

$$\tilde{u}_3(a_3) = e^{-\int_0^a \mu_3(\tau_3)d\tau - (f_{30})_3 a}\left(\tilde{u}_3(0) - w\int_0^a e^{\int_0^\tau \mu_3(s_3)ds + (f_{30})}\tau_3\,(\alpha_3)_3(\tau_3)e^{\gamma_3\tau}d\tau\right).$$

Substituting the expression (13) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\right)\tilde{u}_3(0)$$

$$+w\int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\int_0^a (\alpha_3)_3(\tau_3)e^{-\gamma_3\tau_3}e^{(f_{30})_3\tau_3 + \int_0^\tau \mu_3(s_3)ds}d\tau_3\,da_3$$

$$= 0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(\tau_3)d\tau} da\right)(\tilde{u}_3(0))_3$$

$$+\left(-3-\int_0^\infty e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(\tau_3)d\tau}\int_0^a e^{(f_{30})_3\tau_3+\int_0^\tau \mu_3(s_3)ds+y\tau}\alpha_3(\tau)d\tau\, da\right)w=0. \quad (23)$$

Define $Y(\gamma,a)=e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(t_3)d\tau}$. Then

$$\left(1-\int_0^\infty \beta_3\,(a_3)Y_3(\gamma_3,a_3)da_3\right)u_3(0)$$

$$+w\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)\int_0^a(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\, da_3=0, \quad (24)$$

$$\left(\int_0^\infty Y_3(\gamma_3,a_3)da_3\right)\tilde{u}_3(0)$$

$$+\left(-1-\int_0^\infty Y_3(\gamma_3,a_3)\int_0^a(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\, da_3\right)w=0.$$

Equqtion (20) has a solution of the form of $u_3(t,a_3)=u_3(a_3)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_3(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_3(\gamma_3,a_3)=\int_0^a(\rho_0)_3(\tau_3)Y^{-3}\,(\gamma_3,a_3)d\tau_3$. We obtain the characteristic equation

$$f_{30_3}^l\left(\int_0^\infty Y_3(\gamma_3,a_3)da_3\right)\left(\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$+\left(1-\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)da\right)\left(1+(f_{30}^l)_3\int_0^\infty Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_3(a_3)=(\rho_0)_3(0)e^{-\int_0^a \mu_3(\tau_3)d\tau_3-(f_{30})_3a_3}$$

$$I_3(\gamma_3,a_3)=\int_0^a (\rho_0)_3(\tau_3)e^{\gamma_3\tau_3+(f_{30})_3\tau+\int_0^\tau \mu_3(s_3)ds_3}d\tau$$

$$=\frac{(\rho_0)_3(0)}{\gamma_3}(e^{\gamma_3a_3}-1)$$

Let

$$E_3(a_3)=e^{-(f_{30})_3a_3-\int_0^a \mu_3(\tau_3)d\tau_3}$$

Then $[Y_3]_3[\gamma_3]_3,a_3)=E_3(a_3)e^{-\gamma_3a_3}$. by a simple calculation,

$$\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)I(\gamma_3,a_3)da=\frac{[\rho_0]_3(0)}{\gamma}\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3a_3}\ da\right);$$

$$\int_0^\infty Y_3(\gamma_3,a_3)I(\gamma_3,a_3)da=\frac{[\rho_0]_3(0)}{\gamma_3}\left(\int_0^\infty E_3(a_3)e^{-\gamma_3a_3}\ da_3\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_3(0)}{\gamma_3}[f_{30}^3]\int_0^\infty E_3(a_3)e^{-\gamma_3a_3}da_3\ (\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3a_3}da_3)$$

$$+\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right).$$

$$\left(1+\frac{(\rho_0)_3(0)}{\gamma_3}(f_{30}^{'})_3\left(\int_0^\infty E_3(a_3)da_3-\int_0^\infty E_3(a_3)e^{-\gamma_3 a_3}da_3\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right)(1+\frac{(f_{30}^{'})_3}{\gamma_3}(\rho_0)_3(0)\int_0^\infty E_3(a_3)da_3)=0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_3(a_3)Y_3(\gamma_3 a_3)da_3 = 1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3 \neq 0,$$

Since $\int_0^\infty \beta_3(a_3)Y_3(\gamma_3 a_3)\int_0^a \alpha_3(\tau_3)Y^{-3}(\gamma_3,\tau_3)d\tau_3 da_3 > 0$. Therefore, the characteristic solution is

$$\gamma_3 = -(f_{30}^{'})_3(\rho_0)_3\int_0^\infty E_3(a_3)da_3.$$

Since $f_3'(P_0)_3 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_3), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_3), 0)$ which is related to the extinction of debit summation types

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_3 + (\mu_3(a_3) + [f_3]_3(P_3^*) + [f_2]_3(Q_3^*))P_3^* = 0, \quad (26)$$

$$p_3^*(0) = h_3(Q_3^*)\int_0^\infty \beta_3(a_3)p_3^*(a_J)p_3^*(a_3)da_3 \quad (27)$$

$$g_3(P_3^*) - b_3 = 0, \quad (28)$$

Where $P_3^* = \int_0^\infty p_3^*(a_3)\, da_3$

Equation (28) yields $P^* = g^{-3}(b)$,

From (26),

$$p_3^*(a) = p_3^*(0)e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}$$

Substituting into (27), we need to solve

$$\left(1 - h_3(Q_3^*)\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}da_3\right)p_3^*(0) = 0$$

For $Q_3^*$

Fix $p_3^*(0)$ and $P_3^*$ Let

$$H_3(z_3) = 1 - h_3(z_3)\int_0^\infty \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3} da_{,3}$$

And assume that

$$[k_3]_3 \int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(g_3^{-3}(b_3))a_3 - [k_2]_3 a_3}\, da_3 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_3(z_3)] > 0.$$

Clearly,

$$\frac{dH_3}{dz_3} = h_3(z)([f_2^{'}]_3(z_3))\int_0^\infty a_3\beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - f_3(P_3^*)a_3 - f_2(z_3)a_3} da_{3}$$

$-h^{'}(z_3)\int_0^\infty \beta_3(a_3)\, e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3}\ da_3$

By the assumptions on $\left[f_2^3\right]_3 and\ [h^3]^3 dH_3/dz > 0$. Hence, if

$$H_3(0) = 1 - \int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3} da_3 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_3(0) < 0$ follows if $g_3^-(b_3) < [f_3^{-3}]_3(r_3)$, where $r_3$ is the solution of (12).

Indeed, since

$$R_3(s_3) = \int_0^{\infty} \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 e^{-y\frac{a}{3}}} da_3$$

Is a decreasing function of $s_3$ when $P_3^* = [g^{-3}]_3(b_3) < [f_3^{-3}]_3(r_3)$, that is $[f_3]_3(P_3^*) < r$, $R\big(f_3(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_3(g^{-3}(b))) > 1, i.e.,$

$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3\left(g_3^{-3}(b)\right)a_3} > 1.$ (H4)

Then the system (4)-(6) has a unique positive equilibrium $(P_3^*(a_3), Q_3^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_3^*(a_3), Q_3^*)$ is more complicated. We again consider perturbations $m_3(t_3 a_3) = p_3(t_3, a_3) - p_3^*(a_3)$ and $W_3(t) = Q_3(t) - Q^*$.Then

$M_3(t) \equiv \int_0^{\infty} m_3(t, a_3) da_3 = P_3(t) - P_3^*,$

$$[m_3]_3 + [m_a]_3 = [p_3]_3 + [p_a]_3 - \left[p^*\right]_3(a_3) \approx -(\mu_3(a_3) + [f_{3*}]_3 + [f_{2*}]_3)m_3$$

$$-[f_{3*}^3]_3 p_3^*(a_3)M - [f_{2*}^3]_3 p_3^*(\mathrm{a})W_3,$$

$$m_3(t,0)p_3(t,0) - p_3^*(0) \approx h_* \int_0^{\infty} \beta_3(a_3) m_3(t, a_3)\, da + \frac{[h_*^3]_3}{[h_*]_3}\, p_3^*(0)W_3$$

$$\frac{dW_3}{dt} = \frac{dQ_3}{dt} \approx [g_*^3]_3 Q_3^* M_{,3}$$

Where $[f_{3*}]_3$ denoting $[f_3(P^*)]_3$, $[f_{3*}]_3$ denoting $[f_2(P^*)]_3$, $[f_{3*}^{'}]_3$ denoting $[f^{'}{}_3(P^*)]_3[f_{2*}^{'}]_3$ denoting $[f_2]_3$ $(Q_3^*,\quad h_3^*$ denoting $h_*^{'}$ denoting $h_3^{'}$ $(Q_3^*)$, and $[g_*^{'}]_3$

Denoting $[g^{'}]_3(P_3^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_3]_3+[m_a]_3=-(\mu_3(a_3)+[f_{3*}]_3+[f_{2*}]_3)m_3-[f^{'}{}_{3*}]_3[p_3^*](a_3)M_3-[f_{2*}^{'}]_3p_3^*(a_3)W_3$$

$$m_3(t,0)=[h_*]_3\int_0^\infty \beta_3(a_3)m_3(t,a_3)da_3+\frac{[h_*^3]_3}{[h^*]_3}[p^*]_3(0)\,W_{,3} \qquad (30)$$

$$\frac{dW_3}{dt}=[g_*^3]_3[Q_3^*]M_3$$

Integrating $[m_3]_3(t,a_3)$ with respect to a from 0 to ∞, we have

$$\frac{dM_3(t)}{dt}=\int_0^\infty [mt]_3(t,a_3)da_3$$

$$=-\int_0^\infty [m_a]_3\,(t,a_3)da_3-\int_0^\infty \mu(a_3)m(t,a_3)da_3$$

$$-\left([f_{3*}]_3+[f_{2*}]_3+[f^{'}{}_{3*}]_3\int_0^\infty p_3^*(a_3)da_3\right)M_3-[f_{2*}^3]_3\int_0^\infty p_3^*(a_3)da_3W_3$$

$$=m_3(t,0)-\int_0^\infty \mu_3(a_3)\,m_3(t,a_3)da_3$$

$$-\left([f_{3*}]_3+[f_{2*}]_3+[f_{3*}^{'}]_3P_3^*\right)M_3-[f_{2*}^{'}]_3P_3^*W_{\cdot 3}$$

From (30), we arrive at the equations

$$\frac{dW_3(t)}{dt} = \int_0^{\infty} ([h_*]_3\beta_3(a_3) - \mu_3(a_3))m_3(t,a_3)da_3$$

$$-\left([f_{3*}] + [f_{2*}]_3 + \left[f_{3*}^{'}\right]_3 P_3^*\right) M_3 + \left(\frac{[h_*^{'}]_3}{[h_*]^3}\, P_3^*(0) - f_{2*}^{'} P_3^*\right) W_3,$$

(33)

$$\frac{dW_3(t)}{dt}\left[g^{'}{}_*\right]_3 Q_3^*\, M_{.3}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{"*}]_3\beta_3(a_3) - \mu_3(a_3))m_3(t,a_3)da_3 = ([h_*]_3\beta_3(\theta_3) - \mu_3(\theta_3))M_{,3}$$

Where $0< \theta_3 = \theta_3(t) < \infty$.

And if $a_3^2\,[a_3^2]_3<4\,[a_2]_3<4\,a_2 e$,

$\mathrm{Re}\,\lambda_3\left[\lambda_3\right]_3 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_3 = -\frac{3}{2}a_3[a2\ ]_3 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_*\,[h_*]_3\beta_3(a_3) - \mu_3(a_3)\right|$. Then$|n_3(t)|$ ⊠⊿∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_3^2 <4\,a_2 e$ and $[a_3] > 2\Delta$,or $[a_3^2]$ ⊠$[a_2]e$ and $(a_3 - \sqrt{(a_3^2)_3 - 4(a_2)_3 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+3}\right)$ and $\beta_3(a_3) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3) + d\tau_3 da} = \frac{3}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_3 + [\mu_0]$ 0.5. Assume$[f_3](x_3)= K_3x_3$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_3(x_3)= [0x_3]/(3 + x_3)$; the functional responses $[[f_2](x) = [0]\, x_3/(1 + x_3)$and $h_3(x_3)=(1+\omega_3x_3)/(1 + x_3)$(See Holling [7] and Rashevsky [13])

Equation (34),now has the form of

$$\frac{m_3 + [\mu_0] + [K][P_0])}{[m_3]_3 + (\mu_0)_3 + [K][P_0])_3^3\ [K]_3[P_0]_3^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(3 + P_0) < b_3$i e $[P_0] < b/(\delta_3 - b_3)$then the equilibrium$(P_0)$ $[P_0]_3(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_3$

$[f_{3*}] = K_3\ P^*\ [f_{2*}] = \frac{\sigma Q_3^*}{3+Q_3^*}$ $[h_*] = \frac{3+\omega_3 Q_3^*}{3+Q_3^*}$

$[f_{3*}] = K_3\ [f_{2*}^3] = \frac{\sigma_3}{(3+Q_3^*)^2}[\ g_*^3] = \frac{\sigma_3}{(3+P_3^*)^2}$ $[h_*^3] = \frac{\omega_3-3}{(3+Q_3^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0, K, b, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_3}{(m+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma)^2}+\frac{2\omega_3}{([m_3]_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma_3)^3}<1$$

And

$$\frac{1}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^2}+\frac{2\omega_3}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}+(m_3+[\mu_0]_3+\frac{K_3b_3}{\delta_3-b_3}+\sigma_3)\,Q_3^*)^3$$

$$=m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}(1+Q^*)(1+\omega_3Q_3^*$$

Now

$$\left[a_3^2\right]-4\left[a_2\right]e=\left(\frac{2K_3b_3}{\delta_3-b_3}\right)^2+\frac{4K_3b_3\sigma_3\,Q_3^*}{(\delta_3-b_3)(1+Q_3^*)}+\frac{\sigma_3^2\,Q_3^{*2}}{(1+Q_3^*)^2}$$

$$\frac{4\,b_3\sigma_3\,(\delta_3-b)Q^*}{\delta_3(1+Q_3^*)^2}-\frac{4\,(\delta_3-b_3)^2(1-\omega_3)Q_3^*}{\delta_3(1+\omega_3Q_3^*)(1+Q_3^*)}P*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_3\,ae_3^{mu}\;\;[\mu_0]_3-\frac{a_3}{a_3+3}$$

Hence if $[a_3^2 > 4(a_2)_3 e$ and

$$[a_3] = \frac{2K_3 b_3}{\delta_3 - b_3} + \frac{\sigma_3 Q_3^*}{1 + Q_3^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_3]_3 = [f_{3*}]_3 + [f_{2*}]_3 + [f_{3*}] P_3^*$$

$$[a_2]_3 = \left[f_{2*}^3\right] P_3^* - \frac{h_*^3}{h_*} P^*(0)$$

$$e = [g_*^3]_3 Q_3^*$$

$$\eta_3(t) = \left[h_*\right]_3 \beta_3 \left(\theta_3(t)\right) - \mu_3(\theta(t)),$$

The system(33)can be written as the following system:

$$[M(t)]\left[a_3\right] M_3 - [a_2]_3 W_3 - [n(t)] M_3$$

$$w_3(t) = eM_3$$

Write

$$A = \begin{pmatrix} -[a_3]_3 - [a_2]_3 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_3 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_3 \\ W_3 \end{pmatrix}$$

The solution of (32)is

$$X_3(t) = e^{A,t}[X]_3(0) + \int_0 e^{At}$$

**This equation is useful in following type of problems . For actual model kindly see various papersThere is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant Λ, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive Λ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the Λ>0 generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible.**

**Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of (e) GW signals from BBHs with component spins $\chi_{1,2} \in [-0.99, 0.99]$ aligned (e&eb) with the orbital angular momentum, component masses $m_{1,2} \in [2,48] M_\odot$, and (e&eb) total mass $M_{total} \leq 50\ M_\odot$. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins >+0.6 and component masses $\gtrsim 20\ M_\odot$. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201 STOP Advanced LIGO may be the first experiment to (e) detect gravitational waves. Through superradiance of stellar black holes, it may also be (=) the first experiment to discover the QCD axion with (e&eb) decay constant above the GUT scale. When an axion's Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e) the black hole, forming a "gravitational atom." Through the superradiance process, the number of axions occupying (e) the bound levels grows exponentially, extracting (e) energy and angular momentum from the black hole. Axions**

**transitioning between levels of the gravitational atom and (e&eb) axions annihilating to (e) gravitons can produce (eb) observable gravitational wave signals. The signals are long-lasting, monochromatic, and can be distinguished from (e) ordinary astrophysical sources. Authors estimate up to O (1) transition events at aLIGO for (e) an axion between 10^-11 and (e&eb) 10^-10 eV and up to 10^4 annihilation events for (e) an axion between 10^-13 and 10^-11 eV. In the event of a null search, aLIGO can constrain (e) the axion mass for (e) a range of rapidly spinning black hole formation rates. Axion annihilations are also promising for (e) much lighter masses at future lower-frequency gravitational wave observatories; the rates have (e) large uncertainties, dominated by (e) supermassive black hole spin distributions. Projections for aLIGO are (=) robust against perturbations from (e) the black hole environment and account for (e) updated exclusion on (e&eb) the QCD axion of 6*10^-13 eV < ma < 2*10^-11 eV suggested by (e) stellar black hole spin measurements. Subjects: High Energy Physics - Phenomenology (hep-ph); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 91, 084011 (2015) DOI: 10.1103/PhysRevD.91.084011Cite as:arXiv:1411.2263 [hep-ph] (or arXiv:1411.2263v3 [hep-ph] for this version) Discovering the QCD Axion with Black Holes and Gravitational Waves Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang STOP Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric**

**charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by (e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave**

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**rank-1 CB geometries. Subjects: High Energy Physics - Theory (hep-th) Cite as:arXiv: 1601.00011 [hep-th] (or arXiv: 1601.00011v2 [hep-th] for this version Geometric constraints on the space of N=2 SCFTs II: Construction of special Kähler geometries and RG flows Philip C. Argyres, Matteo Lotito, Yongchao Lü, Mario Martone (Submitted on 31 Dec 2015 (v1), last revised 6 Jan 2016 (this version, v2)) STOP Authors propose a model of Dark Energy in which the field currently dominating (e) the energy density of the universe is (=) an "axion field" linearly coupled to (e&eb) the Pontryagin density, tr (FΛF), (i.e., the exterior derivative of (e) the Chern-Simons form) of (e) a massive gauge field. Authors assume that the axion has (e) self-interactions corresponding to (e&eb) a non-trivial (exponential) potential. Authors argue that a non-vanishing magnetic helicity of (e) the gauge field triggers (eb) slow-rolling of the axion at field values far below (eb) the Planck scale. Proposal leads to (eb) a "Tracking Dark Energy Scenario" in which the contribution of the axion energy density to (e) the total energy density is (=) constant (and small) during (e&eb) the early radiation phase, until (e) a secular growth term proportional to (e) the Pontryagin density of (e) the gauge field becomes (=) dominant. The initially small contribution of (e) the axion field to (e) the total energy density is related to the observed small baryon-to-entropy ratio. Some speculations concerning the nature of the gauge field are offered.**

**Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Cite as:arXiv: 1601.00057 [hep-th] (or arXiv:1601.00057v2 [hep-th] for this version) Tracking Dark Energy from Axion-Gauge Field Couplings Stephon Alexander, Robert Brandenberger, Juerg Froehlich (Submitted on 1 Jan 2016 (v1), last revised 23 Jan 2016 (this version, v2)) STOP Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a**

**pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version) Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP In this work authors obtain bounds on (e&eb) the topological Abelian string-vortex and on (e&eb) the string-cigar, by using (e) a new measure of configurational complexity, known as (=) configurational entropy. In this way, the information-theoretical measure of (e) six-dimensional braneworld scenarios is capable to (e) probe situations where (e) the parameters responsible for (e) the brane thickness are arbitrary. The so-called configurational entropy (CE) selects (e&eb) the best value of the parameter in (eb) the model. This is accomplished by (e) minimizing the CE, namely, by (e) selecting the most appropriate parameters in the model that correspond to (e&eb) the most organized system, based upon (e) the Shannon information theory. This information-theoretical measure of complexity provides (eb) a complementary perspective to (e) situations where (e) strictly energy-based arguments are (=) inconclusive. Authors show that the higher the energy the higher the CE, what shows (eb) an important correlation between the energy of the localized field configuration and (e&eb) its**

**associated entropic measure. Subjects: High Energy Physics - Theory (hep-th) DOI:10.1016/j.physletb.2016.02.038 Cite as: arXiv: 1601.00076 [hep-th] (or arXiv: 1601.00076v2 [hep-th] for this version) Bounds on topological Abelian string-vortex and string-cigar from information-entropic measure R. A. C. Correa, D. M. Dantas, C. A. S. Almeida, Roldao da Rocha (Submitted on 1 Jan 2016 (v1), last revised 17 Feb 2016 (this version, v2)) STOP Sum of the Holevo quantity (that bounds the capacity of quantum channels to transmit (e&eb) classical information about an observable) and (e&eb) the quantum discord (a measure of the quantumness of (e) correlations of that observable) yields (eb) an observable-independent total given by (e) the quantum mutual information This split naturally delineates ,dilates and disseminates information about (e&eb) quantum systems accessible to (e) observers -- information that is redundantly transmitted by (e&eb) the environment -- while showing that it is (=) maximized for (e) the quasi-classical pointer observable. Other observables are accessible only via (e&eb) correlations with the pointer observable .Further authors prove an anti-symmetry property relating accessible information and (e&eb) discord. It shows that information becomes (=) objective -- accessible to (e) many observers -- only as quantum information is relegated to (e) correlations with (e&eb) the global environment, (eb) and, therefore, locally inaccessible. The resulting complementarity explains (eb) why, in a quantum Universe, we perceive objective classical reality while (e) flagrantly quantum superpositions are out of reach. Subjects: Quantum Physics (quant-ph) Journal reference: Scientific Reports 3, 1729 (2013) Report number: LAUR-13-26386 Cite as:arXiv: 1303.4659 [quant-ph] (or arXiv: 1303.4659v2 [quant-ph] for this version) Complementarity of quantum discord and classically accessible information Michael Zwolak, Wojciech H. Zurek STOP Niels Bohr proposed that the outcome of the measurement becomes (=) objective and real, and, hence, classical, when (e) its results can be communicated by (e) classical means. In this work authors revisit Bohr's postulate using (e) modern tools from (e) the quantum information theory. They find a**

**full confirmation of Bohr's idea: if a measurement device is in a nonclassical state, (eb) the measurement results cannot be communicated perfectly by classical means. In this case some part of information in the measurement apparatus is lost (e) in the process of communication: the amount of this lost information turns out to be (=) the quantum discord. The information loss occurs even when (e) the apparatus is not entangled with the system of interest. The tools presented in this work allow (eb) to generalize Bohr's postulate: authors show that for pure system-apparatus states (eb) quantum communication does not provide any advantage when (e) measurement results are communicated to (e) more than one recipient. They further demonstrate the superiority of quantum communication to (e&eb) two recipients on (e&eb) a mixed system-apparatus state and show that (eb) this effect is fundamentally different from (e) quantum state cloning. Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. Lett. 111, 040401 (2013) DOI:10.1103/PhysRevLett.111.040401 Report number: LA-UR-12-25351 Cite as: arXiv: 1304.2914 [quant-ph] (or arXiv: 1304.2914v2 [quant-ph] for this version) Quantum discord cannot be shared Alexander Streltsov, Wojciech H. Zurek STOP Wojciech H. Zurek, Don N. Page investigate thermodynamic equilibrium of (e) a self-gravitating perfect fluid in (eb) a spherically symmetric system containing (e) a black hole of mass M by means of (e) the Tolman-Oppenheimer-Volkoff (TOV) equation. At r >> 2M its solutions describe (eb) a black-body radiation atmosphere with (e&eb) the Hawking temperature T_BH~1/(8 \pi M) that is increasingly blueshifted as r approaches (eb) 2M. However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments:**

**Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by**

**(e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a firewall-like "Planck cocoon" with entropy close to (e) black hole entropy**

# Prescientiаl Dynamics And Contential Staticity Of Accentuation Attrition Models

## Module Four

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equailibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-445) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Some of the concepts to which model is applicable is as follows:**

**Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by (e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave signal. The signal sweeps upwards in frequency from 35 to 250 Hz with (e&eb) a peak gravitational-wave strain of (e) 1.0×10−21. It matches the waveform predicted by general relativity for (e) the inspiral and merger of a pair of black holes and the ringdown of the resulting (eb) single black hole. The signal was observed with a**

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Almeida, Roldao da Rocha (Submitted on 1 Jan 2016 (v1), last revised 17 Feb 2016 (this version, v2)) STOP Sum of the Holevo quantity (that bounds the capacity of quantum channels to transmit (e&eb) classical information about an observable) and (e&eb) the quantum discord (a measure of the quantumness of (e) correlations of that observable) yields (eb) an observable-independent total given by (e) the quantum mutual information This split naturally delineates ,dilates and disseminates information about (e&eb) quantum systems accessible to (e) observers -- information that is redundantly transmitted by (e&eb) the environment -- while showing that it is (=) maximized for (e) the quasi-classical pointer observable. Other observables are accessible only via (e&eb) correlations with the pointer observable .Further authors prove an anti-symmetry property relating accessible information and (e&eb) discord. It shows that information becomes (=) objective -- accessible to (e) many observers -- only as quantum information is relegated to (e) correlations with (e&eb) the global environment, (eb) and, therefore, locally inaccessible. The resulting complementarity explains (eb) why, in a quantum Universe, we perceive objective classical reality while (e) flagrantly quantum superpositions are out of reach. Subjects: Quantum Physics (quant-ph) Journal reference: Scientific Reports 3, 1729 (2013) Report number: LAUR-13-26386 Cite as:arXiv: 1303.4659 [quant-ph] (or arXiv: 1303.4659v2 [quant-ph] for this version) Complementarity of quantum discord and classically accessible information Michael Zwolak, Wojciech H. Zurek STOP Niels Bohr proposed that the outcome of the measurement becomes (=) objective and real, and, hence, classical, when (e) its results can be communicated by (e) classical means. In this work authors revisit Bohr's postulate using (e) modern tools from (e) the quantum information theory. They find a full confirmation of Bohr's idea: if a measurement device is in a nonclassical state, (eb) the measurement results cannot be communicated perfectly by classical means. In this case some part of information in the measurement apparatus is lost (e) in the process of communication: the amount of this lost information turns out to be (=) the quantum discord. The information loss occurs even when (e) the apparatus is not entangled with the system of interest. The tools presented in this work allow (eb) to generalize Bohr's postulate: authors show that for pure system-apparatus states (eb) quantum communication does not provide any advantage when (e) measurement results are communicated to (e) more than one recipient. They further demonstrate the superiority of quantum communication to (e&eb) two recipients on (e&eb) a mixed system-apparatus state and show that (eb) this effect is fundamentally different from (e) quantum state cloning. Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. Lett. 111, 040401 (2013) DOI:10.1103/PhysRevLett.111.040401 Report number: LA-UR-12-25351 Cite as: arXiv: 1304.2914 [quant-ph] (or arXiv: 1304.2914v2 [quant-ph] for this version) Quantum discord cannot be shared Alexander Streltsov, Wojciech H. Zurek STOP Wojciech H. Zurek, Don N. Page investigate thermodynamic equilibrium of (e) a self-gravitating perfect fluid in (eb) a spherically symmetric system containing (e) a black hole of mass M by means of (e) the Tolman-Oppenheimer-Volkoff (TOV) equation. At r >> 2M its solutions describe (eb) a black-body radiation atmosphere with (e&eb) the Hawking temperature T_BH~1/(8 \pi M) that is increasingly blueshifted as r approaches (eb) 2M. However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments: Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by (e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a**

**firewall-like "Planck cocoon" with entropy close to (e) black hole entropy. Authors reproduce paper with an updated title and abstract. For a later study with similar results see G. 't Hooft, gr-qc/9706058 Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Physical Review D, Volume 29, Number 4, pp. 628 - 631 (1984) DOI:10.1103/PhysRevD.29.628 Cite as:arXiv:1511.07051 [gr-qc] (or arXiv:1511.07051v1 [gr-qc] for this version) Firewalls, black-hole thermodynamics, and singular solutions of the Tolman-Oppenheimer-Volkoff equation Wojciech H. Zurek, Don N. Page (Submitted on 22 Nov 2015) STOP Frederic Magniez, Ashwin Nayak consider the problem of testing the commutativity of (e) a black-box group specified by its k generators. The complexity (in terms of k) of (e) this problem was first considered by Pak, who gave (eb) a randomized algorithm involving O(k) group operations. They construct a quite optimal quantum algorithm for this problem whose complexity is in O~(k2/3)O~(k2/3). The algorithm uses and highlights the power of the quantization method of Szegedy. For the lower bound of Ω(k2/3)Ω(k2/3), authors give a reduction from (e) a special case of Element Distinctness to our problem. Along the way, they prove the optimality of the algorithm of (e) Pak for the randomized model. Algorithmica July 2007, Volume 48, Issue 3, pp 221-232 First online: 25 June 2007 Quantum Complexity of Testing Group Commutativity Frederic Magniez, Ashwin Nayak STOP Pawel Wocjan, Jon Yard analyze relationships between quantum computation and (e&eb) a family of generalizations of the Jones polynomial. Extending recent work by Aharonov et al., authors give efficient quantum circuits for (e) implementing the unitary Jones-Wenzl representations of the braid group. We use these to provide new quantum algorithms for (e) approximately evaluating a family of specializations of the HOMFLYPT two-variable polynomial of trace closures of braids. They also give (eb) algorithms for approximating the Jones polynomial of a general class of closures of (e) braids at roots of unity. Next authors provide (eb) a self-contained proof of a result of (e) Freedman et al. that any quantum computation can be replaced by (e&eb) an additive approximation of the Jones polynomial, evaluated at (eb) almost any primitive root of unity. Proof encodes two-qubit unitaries into (e&eb) the rectangular representation of (e) the eight-strand braid group. Authors then give QCMA-complete and (e&eb) PSPACE-complete problems which are based on braids. They conclude with direct proofs that evaluating the Jones polynomial of the plat closure at most primitive roots of unity is a #P-hard problem, while learning its most significant bit is PP-hard, circumventing (e) the usual route through the Tutte polynomial and graph coloring. Subjects: Quantum Physics (quant-ph); Geometric Topology (math.GT) Cite as: arXiv:quant-ph/0603069 (or arXiv:quant-ph/0603069v3 for this version) The Jones polynomial: quantum algorithms and applications in quantum complexity theory Pawel Wocjan, Jon Yard STOP Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan address the task of switching off the Hamiltonian of (e) a system by removing (e) all internal and system-environment couplings. Authors propose dynamical decoupling schemes, which use (e) only bounded-strength controls, for quantum many-body systems with local system Hamiltonians and local environmental couplings. To do so, they introduce (eb) the combinatorial concept of (e) balanced-cycle orthogonal arrays (BOAs) and show how to construct them from (e) classical error-correcting codes. The derived decoupling schemes may be useful as (=) a primitive for more complex schemes, e.g., for Hamiltonian simulation. For the case of n qubits and a 2-local Hamiltonian, the length of the resulting decoupling (e&eb) scheme scales as O (nlogn), improving over (e&eb) the previously best-known schemes that scaled quadratically with (e&eb) n. More generally, using balanced-cycle orthogonal arrays constructed from families of BCH codes, we show that bounded-strength decoupling for any ℓ-local Hamiltonian, where ℓ≥2, can be achieved using decoupling schemes of length at most O(nℓ−1logn). Subjects:Quantum Physics (quant-ph); Emerging Technologies (cs.ET) Journal reference: IEEE Trans. Inf. Theory, 62(5): 2881-2894, 2016 DOI: 10.1109/TIT.2016.2535183 Report number: MIT-CTP /4705 Cite as: arXiv: 1509.00408 [quant-ph] (or arXiv: 1509.00408v2 [quant-ph] for this version) Improved bounded-strength decoupling schemes for local Hamiltonians Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan STOP Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola propose dynamical control schemes for Hamiltonian simulation in many-body quantum systems that avoid instantaneous control operations and rely solely on (e) realistic bounded-strength control Hamiltonians. Each simulation protocol consists of (e) periodic repetitions of a basic control block, constructed as (=) a suitable modification of an "Eulerian decoupling cycle," that would otherwise implement (eb) a trivial (zero) target Hamiltonian. For an open quantum system coupled to (e) an uncontrollable environment, approach may be employed to (e) engineer an effective evolution that simulates (e&eb) a target Hamiltonian on the system, while suppressing (e) unwanted decoherence to the leading order. Authors present illustrative applications to (e) both closed- and open-system simulation settings, with (e&eb)**

**emphasis on simulation of non-local (two-body) Hamiltonians using (e) only local (one-body) controls. In particular, authors provide (eb) simulation schemes applicable to (e&eb) Heisenberg-coupled spin chains exposed to (e) general linear decoherence, and show how to simulate (e&eb) Kitaev's honeycomb lattice Hamiltonian starting from Ising-coupled qubits, as (=) potentially relevant to (e&eb) the dynamical generation of a topologically protected (e&eb) quantum memory. Additional implications for quantum information processing are discussed. Subjects: Quantum Physics (quant-ph) Journal reference: New J. Phys. 16, 045021 (2014) DOI: 10.1088/1367-2630/16/4/045021 Cite as:arXiv: 1310.4153 [quant-ph]**

**(or arXiv: 1310.4153v1 [quant-ph] for this version) Hamiltonian quantum simulation with bounded-strength controls Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola STOP Felix Fontein, Pawel Wocjan present a quantum algorithm for computing the period lattice of (e) infrastructures of fixed dimension. The algorithm applies to (e&eb) infrastructures that satisfy certain conditions. The latter are always fulfilled for (e) infrastructures obtained from (e) global fields, i.e., algebraic number fields and function fields with (e&eb) finite constant fields. The first main contributions is an exponentially better method for (e) sampling approximations of vectors of the dual lattice of (e) the period lattice than the methods outlined in the works of (e) Hallgren and Schmidt and Vollmer. This new method improves (eb) the success probability by (e) a factor of at least 2^ {n^2-1} where (e) n is the dimension. The second main contribution is a rigorous and complete proof that the running time of the algorithm is (=) polynomial in the logarithm of (e) the determinant of the period lattice and (e&eb) exponential in n. The third contribution is the determination of (e) an explicit lower bound on (e&eb) the success probability of our algorithm which greatly improves on (e&eb) the bounds given in the above works. The exponential scaling seems inevitable because (e) the best currently known methods for (e) carrying out fundamental arithmetic operations in infrastructures obtained from (e) algebraic number fields take exponential time. In contrast, the problem of computing (e&eb) the period lattice of infrastructures arising from function fields can be solved without (e) the exponential dependence on (e&eb) the dimension n since this problem reduces efficiently to the abelian hidden subgroup problem. This is also true for (e) other important computational problems in (eb) algebraic geometry. The running time of the best classical algorithms for (e) infrastructures arising from (e) global fields increases subexponentially with (e&eb) the determinant of the period lattice. Subjects: Quantum Physics (quant-ph) Cite as: arXiv: 1111.1348 [quant-ph]**

**(or arXiv: 1111.1348v2 [quant-ph] for this version) Quantum Algorithm for Computing the Period Lattice of an Infrastructure Felix Fontein, Pawel Wocjan STOP Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan advocate a new approach of (e) addressing hidden structure problems and (e&eb) finding efficient quantum algorithms. Authors introduce and investigate the Hidden Symmetry Subgroup Problem (HSSP), which is a generalization of the well-studied Hidden Subgroup Problem (HSP). Given a group acting on (e&eb) a set and an oracle whose level sets define (eb) a partition of the set; the task is to recover (e) the subgroup of symmetries of (e) this partition inside the group. The HSSP provides (eb) a unifying framework that, besides the HSP, encompasses (e) a wide range of algebraic oracle problems, including (e) quadratic hidden polynomial problems. While the HSSP can have (e) provably exponential quantum query complexity, authors obtain efficient quantum algorithms for various interesting cases. To achieve this, authors present a general method for (e) reducing the HSSP to the HSP, which works efficiently in several cases related to (e&eb) symmetries of polynomials. The HSSP therefore connects in a rather surprising way certain hidden polynomial problems with (e&eb) the HSP. Using this connection, authors obtain the first efficient quantum algorithm for (e) the hidden polynomial problem for (e) multivariate quadratic polynomials over fields of (e) constant characteristic. Authors also apply the new methods to (e&eb) polynomial function graph problems and (e&eb) present an efficient quantum procedure for (e) constant degree multivariate polynomials over (e&eb) any field. This result improves in several ways the currently known algorithms. Subjects: Quantum Physics (quant-ph) Journal reference: SIAM Journal on Computing 42:5 (2013) DOI: 10.1137/120864416 Cite as: arXiv: 1107.2189 [quant-ph]**

**(or arXiv: 1107.2189v2 [quant-ph] for this version) Hidden Symmetry Subgroup Problems Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan STOP Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan present an efficient general method for realizing a quantum walk operator corresponding to (e&eb) an arbitrary sparse classical random walk. Approach is based on Grover and Rudolph's method for preparing coherent versions of (e) efficiently integrable probability distributions. This method is intended for (e) use in quantum walk algorithms with polynomial speedups, whose complexity is usually measured in terms of (e&eb) how many times we have to apply a step of a quantum walk, compared to (e&eb) the number of necessary classical Markov chain steps. Authors consider a finer notion of complexity including (e) the**

**number of elementary gates it takes to implement each step of the quantum walk with (e&eb) some desired accuracy. The difference in complexity for (e) various implementation approaches is that method scales linearly in the sparsity parameter and (e&eb) poly-logarithmically with (e&eb) the inverse of the desired precision. Best previously known general methods either scale quadratically in (eb) the sparsity parameter, or polynomially in (eb) the inverse precision. Approach is especially relevant for (e) implementing quantum walks corresponding to (e&eb) classical random walks like those used in (eb) the classical algorithms for approximating permanents and sampling from binary contingency tables. In those algorithms, the sparsity parameter grows with (e&eb) the problem size, while (e) maintaining high precision is required. Subjects: Quantum Physics (quant-ph) Journal reference: Quantum Information and Computation, Vol.10, No.5&6, pp0420-0434 (2010) Cite as: arXiv: 0903.3465 [quant-ph] (or arXiv: 0903.3465v3 [quant-ph] for this version) Efficient Circuits for Quantum Walks Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan STOP Preparing the ground state of a system of interacting classical particles is (=) an NP-hard problem. Thus, there is in general no better algorithm to (e) solve this problem than exhaustively going through all N configurations of (e) the system to determine (eb) the one with lowest energy, requiring (e) a running time proportional to N. A quantum computer, if it could be built, could solve (e) this problem in time sqrt (N). Here, authors present a powerful extension of this result to (e) the case of interacting quantum particles, demonstrating (eb)that a quantum computer can prepare the ground state of (e) a quantum system as efficiently as it does for (e) classical systems. Subjects: Quantum Physics (quant-ph) Journal reference: Phys.Rev.Lett.102:130503, 2009 DOI:10.1103/PhysRevLett.102.130503 Cite as: arXiv: 0809.2705 [quant-ph] (or arXiv: 0809.2705v1 [quant-ph] for this version) Preparing ground states of quantum many-body systems on a quantum computer David Poulin, Pawel Wocjan STOP Daniel Nagaj, Pawel Wocjan construct a simple translationally invariant, nearest-neighbor Hamiltonian on (e&eb) a chain of 10-dimensional qudits that makes it possible to (e) realize universal quantum computing without (e) any external control during (e&eb) the computational process. Authors only require the ability to prepare an initial computational basis state which encodes (e) both the quantum circuit and its input. The computational process is (=) then carried out by (e) the autonomous Hamiltonian time evolution. After a time polynomially long in the size of (e) the quantum circuit has passed, the result of (e) the computation is obtained with (e&eb) high probability by (e) measuring a few qudits in the computational basis. This result also implies (eb) that there cannot (e) exist efficient classical simulation methods for (e) generic translationally invariant nearest-neighbor Hamiltonians on (e&eb) qudit chains, unless (e) quantum computers can be efficiently simulated by (e&eb) classical computers (or, put in complexity theoretic terms, unless BPP=BQP). Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. A 78, 032311 (2008) DOI: 10.1103/PhysRevA.78.032311 Cite as: arXiv: 0802.0886 [quant-ph] (or arXiv: 0802.0886v4 [quant-ph] for this version) Hamiltonian Quantum Cellular Automata in 1D Daniel Nagaj, Pawel Wocjan**

## 37. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_4}{dt} = -b_4 Q_4 + g_4(g_4)Q_4$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_4(0) = 0, \quad g_4(P_4) > 0, \; P_4 \; 0, \qquad \lim_{P_4 \to \infty} g_4(P_4) = (K_4)_4 \; \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_4 + (\rho_a)_4 + D_4(a_4, P_4, Q_4)\rho_4 + 0, \qquad 0 < a_4 \; \infty,$$

$$\rho_4(t, 0) = \int_0^\infty n_4(a_4, P_4, Q_4)\rho_4(t_4, a_4)da_4$$

Where $D_4$ is the mortality function and $n_4$ is the fertility function.

Put $_4$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_4}{dt}\right]_4 = x_4\left[\frac{dQ_2}{dt}\right]_4 - x_2\left[\frac{dQ_4}{dt}\right]_4$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_4 - x_4\left[\frac{dQ_2}{dt}\right]_4$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_4 - x_6\left[\frac{dQ_3}{dt}\right]_4$$

Where $x_4, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_4$ and $g_4(P_4)Q_4$

Now when you add the term $(b_4Q_4)$ with $g_4(P_4)(Q_4)$ we are infect see is that $(b_4Q_4)$ is credited incrementally part of $g_4(P_4)(Q_4)$ or instant nearly, when subtracted from $\frac{dQ_4}{dt}$, it means that $g_4(P_4)(Q_4)$ is slowly eaten by $\frac{dQ_4}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_4 = \mu_4(a_4) + (f_4)_4(P_4)_4 + (f_2)_4(Q)_4 \quad (2)$$

$$n_4 = \beta_4(a_4)h_4(Q_4), \quad (3)$$

Where $\mu_4$the age- dependent debit rate whit is $\mu_4(a_4)$ $\forall a > 0$; $(f_4)_4$ is the density- dependent death rate with $f_4(0) = 0$, $(f_4^4)_4(P_4)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_4)_4(P_4) = (k_4)_4$ ; $f_2$ is the functional response of creditors to prey with $[f_2]_4(0) = 0, (f_2^4)_4(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_4(Q_4) = [k_2]$ ; $\beta$ is the age dependent birth rate with $\beta(a)$ $0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 4, h^4(Q)$ $0, \forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_4 + (\rho_a)_4 = -\mu_4(a_4) + (f_4)_4P_4) + \big(f_2)_4(Q_4)\big)(\rho_4)_4 \quad (4)$$

$$\rho_4(t_4, 0) = h_4(Q_4)\int_0^\infty \beta_4(a_4)\rho_4(t, a_4)da, \quad (5)$$

$$\frac{dQ_4}{dt}$$

$$= -b_4Q_4 + g_4(P_4)Q_4 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_4$ and $Q_4$ of (4-6) satisfy

$$0 \leq \rho_4(t, a_4) \leq \rho_4^*(t, a_4),$$

$$0 \leq Q_4(t) \leq Q_4^*(t),$$

Where $\rho_4^*(t, a_4)$ and $Q_4^*(t)$, are solutions of

$$(\rho_t^*)_4 + (\rho_a^*)_4 = -(\mu_4(a_4) + (f_4)_4(P_4^*) + (f_2)_4 Q_4^*))\, \rho_4^*, \quad (7)$$

$$\rho_4^*(t, 0) = \int_0^\infty \beta_4(a_4)\, \rho_4^*(t, a_4) da_4 \quad (8)$$

$$\frac{dQ_4^*}{dt} = -b_4 Q_4^* + g_4(P_4^*)Q_4^*, \quad (9)$$

With $P_4^*(t) = \int_0^\infty \rho_4^*(t, a_4) da_4$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_4^*(t)}{dt} = P_4^*(t)\big(r_4 - (f_4)(P_4^*) - (f_2)_4(Q^*)\big), \quad (10)$$

$$\frac{dQ_4^*(t)}{dt} = Q_4^*(t)(-b_4 + g_4(P_4^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_4(a) e^{-\int_0^a \mu_4(\tau)dt} e^{-ra}\, da_4 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_4^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^{\infty} \beta_4(a_4) e^{-\int_0^a \mu_4(t)dt}\, da_4 > 1, \qquad (H1)$$

Which guarantees $r_4 > 0$.

*Remark.* If $\beta_4(a) = [\beta_0]_4 a_4 e^{-\alpha_4 a_4}, \mu_4(a_4) \equiv \mu_4$ then the solution of (12) is

$$r_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_4]_4(P_4) = 0, [f_2](Q_4) = \lambda_4 Q_4\ h_4(Q_4) = 1, \qquad and\ \ g_4(P_4) = c_4 P_4$$

$$[\rho_4]_4 + [\rho_a]_4 = -(\mu_4 + \lambda_4 Q_4)\rho_4$$

$$\rho_4(t,0) = \int_0^{\infty} [\beta_0]_4 a_4 e^{-\alpha_4 a_4} \rho_4(t, a_4)\, da_4$$

$$\frac{dQ_4}{dt} = -b_4 Q_4 + c_4 P_4 Q_4$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_4}{dt} = \omega_4 P_4 - \lambda_4 P_4 Q_4$$

$$\frac{dQ_4}{dt} = b_4 Q_4 - c_4 P_4 Q_4$$

Where $\omega_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$

## 38. THE EQUILIBRIUM $(\rho_0)_4(a_4), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_4(a_4), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_4 + (\rho_a)_4 = -\big(\mu_4(a_4) + (f_4)_4(P_4)\big)\rho,$$

$$\rho_4(t,0) = \int_0^\infty \beta_4(a_4)\rho_4(t,a_4)da_4 \qquad (13)$$

The solution of (13) has the form of

$$\rho_4(t,a_4) = Ce_4^{r(t-a_4)} - \int_0^a (\tau)d\tau - \int_0^4 (f_4)_4\big(P_4(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_4(t) = C_4 e^{r_4 t - \int_0^t (f_4)_4(P_4(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_4^{-4}(r)$ if $k_4 > r$. This is equivalent to

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{(k_4)_4 a_4} da_4 < 1.$$

In fact, if we define

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau)d\tau} e^{-s_4 a_4}\, da_4$$

$R(s)$ is a decreasing function. $R_4(r_4) = 1$ And $R_4(k_4)_4 < 1$ imply $r_4 < (k_4)_4$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_4(a_4), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{k_4 a_4} da_4$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_4 \equiv \int_0^\infty (\rho_0)_{44}(a_4)da_4$ *satisfies*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{-(f_4)_4(P_0)a_4} da$$

$$= 1 \qquad (11)$$

Next, we study the stability of $(\rho_0)_4(a_4), 0)$

Let $u_4(t, a_4) = \rho(t, a_4) - \rho_0(a_4)$ Then

$$N_4(t) \equiv \int_0^\infty u_4(t, a_4)da_4 = \int_0^\infty \rho_4(t, a_4)da_4 - \int_0^\infty (\rho_0)_4(a_4)da = P_4(t) - P_0)_4$$

And

$$(u_4)_4 + (u_a)_4 = (\rho_t)_4 + (\rho_a)_4 - (\rho_0^4)_4(a_4)$$

$$= -\mu_4(a_4) + (f_4)_4(P_4) + (f_2)_4(Q_4))\,(u_4 + (\rho_0)_4) + \mu_4(a_4)(\rho_0)_4 + (f_4)_4(P_0)_4(\rho_0)_4$$

$$= -\mu_4(a_4)u_4 - (f_4)_4(P_4) + (f_2)_4(Q_4))\,u_4 - (\rho_0)_4\,(f_4)_4(P_4) - (f_4)_4(P_0)_4$$

$$+ (f_2)_4(Q_4)), \qquad (15)$$

Where $(\rho_0^4)_4(a_4)$ denotes $d(\rho_0)_4(a_4)/da_4$

Expansion $(f_4)_4(P_4)$ and $(f_2)_4(Q_4)$ at $(P_0)_4$ and 0 yield respectively

$$(f_4)_4(P_4) = (f_4)_4((P_0)_4 + N_4)$$

$$= (f_4)_4(P_0)_4 + (f_4^4)_4(P_0)_4 N_4 + o_4(N_4), \quad (16)$$

$$(f_2)_4(Q_4)$$

$$= (f_2^4)_4(0)Q + o_4(Q_4) \quad (17)$$

Write $f_4(P_0) = f_{40}, f_4'(P_0) = f_{40}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_4)_4 + (u_a)_4 \approx -(\mu_4(a_4) + (f_{40})_4)u - (f_{40}^4)_4(\rho_0)_4(N_4) - (f_{20}^4)_4(\rho_0)_4(Q_4)$$

$$u_4(t,0) = \rho_4(t,0)(\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a) + (\rho_0)_4(a_4)da_4 - (\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a_4)da_4 + (h_4(Q_4) - 4)(\rho_0)_4(0) \quad (18)$$

Expanding

$$h_4(Q_4) = 1 + (h^4)_4(0)Q_4 + 0(Q_4) \equiv 1 + h_0'Q_4 + 0Q_4$$

And substituting it into (48), we arrive at

$$u_4(t,0) \approx \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + (h_{a_4}')(\rho_0)_4 Q_4$$

The linearization is the following :

$$[u_4] + [u_4] = -(\mu_4(a_4) + [t_{40}]_4 u_4 - [\alpha_4]_4(a_4)N_4 - [\alpha_2]_4(a_4)Q_4$$

$$u_4(t,0) = \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + [\xi_4]Q_4 \quad (19)$$

$$\frac{dQ_4}{dt} = -b_4 Q_4$$

where

$$[\alpha_4]_4(a_4) = [t_4^4 a]_4[p, o]_4(a_4) d2(a_4) = [f_{20}^4]_4(\rho\, o\,)_4(a_4)\xi = [\mathrm{h\ o}^4]_4[\rho \mathrm{o}]_4(\mathrm{o}) \text{ and } \mathrm{b}_4 = \mathrm{b}_4 - g_4(P_4)$$

If $g_4(P_0) > b$, clearly, $(\rho_0)_4(a_4), 0)$ is unstable and the predator population is persistent. Assume $g_4(P_0)_4 < b$.

Then, $Q_4(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_4 + u_a = -(\mu_4(a_4) + (f_{40})_4)u_4 - (\alpha_4)_4(a_4)N_4$$

$$= -(\mu_4(a_4) + f_{40})u_4 - \alpha_4(a_4)\int_0^\infty u_4(t, a_4)da_4$$

$$u_4(t, 0) = \int_0^\infty \beta_4\,(a_4)u_4(t, a_4)da_4$$

Using the method developed by Cushing [2, 3], we suppose

$$u_4(t, a_4) = \tilde{\mathrm{u}}_4(a_4)^{e^{\gamma(t-a_4)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_4}\, da_4$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_4)$ gives,

$$\tilde{\mathrm{u}}_4(a_4) = e^{-\int_0^a \mu_4(\tau_4)d\tau - (f_{40})_4 a}\left(\tilde{\mathrm{u}}_4(0) - w\int_0^a e^{\int_0^\tau \mu_4(s_4)ds + (f_{40})}\tau_4\,(\alpha_4)_4(\tau_4)e^{\gamma_4\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1-\int_0^{\infty}\beta_4(a_4)e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\right)\tilde{u}_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a(\alpha_4)_4(\tau_4)e^{-\gamma_4\tau_4}e^{(f_{40})_4\tau_4+\int_0^\tau \mu_4(s_4)ds}d\tau_4\,da_4$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}da\right)(\tilde{u}_4(0))_4$$

$$+\left(-4-\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a e^{(f_{40})_4\tau_4+\int_0^\tau \mu_4(s_4)ds+y\tau}\alpha_4(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(t_4)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_4\,(a_4)Y_4(\gamma_4,a_4)da_4\right)u_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)\int_0^a(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4=0,$$

(24)

$$\left(\int_0^{\infty}Y_4(\gamma_4,a_4)da_4\right)\tilde{u}_4(0)$$

$$+\left(-1-\int_0^{\infty} Y_4(\gamma_4,a_4)\int_0^{a}(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4\right)w=0.$$

Equqtion (20) has a solution of the form of $u_4(t,a_4)=u_4(a_4)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_4(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_4(\gamma_4,a_4)=\int_0^a(\rho_0)_4(\tau_4)Y^{-4}(\gamma_4,a_4)d\tau_4$. We obtain the characteristic equation

$$f_{40}^{l}{}_4\left(\int_0^{\infty} Y_4(\gamma_4,a_4)da_4\right)\left(\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$+\left(1-\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)da\right)\left(1+(f_{40}^{l})_4\int_0^{\infty}Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_4(a_4)=(\rho_0)_4(0)e^{-\int_0^a\mu_4(\tau_4)d\tau_4-(f_{40})_4a_4}$$

$$I_4(\gamma_4,a_4)=\int_0^a(\rho_0)_4(\tau_4)e^{\gamma_4\tau_4+(f_{40})_4\tau+\int_0^{\tau}\mu_4(s_4)ds_4}d\tau$$

$$=\frac{(\rho_0)_4(0)}{\gamma_4}(e^{\gamma_4a_4}-1)$$

Let

$$E_4(a_4)=e^{-(f_{40})_4a_4-\int_0^a\mu_4(\tau_4)d\tau_4}$$

Then $[Y_4]_4[\gamma_4]_4, a_4) = E_4(a_4)e^{-\gamma_4 a_4}$. by a simple calculation,

$$\int_0^\infty \beta_4(a_4)Y_4(\gamma_4, a_4)I(\gamma_4, a_4)da = \frac{[\rho_0]_4(0)}{\gamma}\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}\ da\right);$$

$$\int_0^\infty Y_4(\gamma_4, a_4)I(\gamma_4, a_4)da = \frac{[\rho_0]_4(0)}{\gamma_4}\left(\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}\ da_4\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_4(0)}{\gamma_4}[f_{40}^4]\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}da_4\ (\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4)$$

$$+\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right).$$

$$\left(1 + \frac{(\rho_0)_4(0)}{\gamma_4}(f_{40}^{'})_4\left(\int_0^\infty E_4(\,a_4)da_4 - \int_0^\infty E_4(\,a_4)e^{-\gamma_4 a_4}da_4\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right)(1 + \frac{(f_{40}^{'})_4}{\gamma_4}(\rho_0)_4(0)\int_0^\infty E_4(\,a_4)da_4) = 0.$$

We need ũ(a) $\not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1 - \int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)da_4 = 1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4 \neq 0,$$

Since $\int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)\ \int_0^a \alpha_4\,(\tau_4)Y^{-4}(\gamma_4, \tau_4)d\tau_4 da_4 > 0$. Therefore, the characteristic solution is

$$\gamma_4 = -(f_{40}^{'})_4(\rho_0)_4 \int_0^{\infty} E_4(a_4)da_4.$$

Since $f_4'(P_0)_4 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_4), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_4), 0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_4 + (\mu_4(a_4) + [f_4]_4(P_4^*) + [f_2]_4(Q_4^*))P_4^* = 0, \quad (26)$$

$$p_4^*(0) = h_4(Q_4^*)\int_0^{\infty} \beta_4(a_4)p_4^*(a_J)p_4^*(a_4)da_4 \quad (27)$$

$$g_4(P_4^*) - b_4 = 0, \quad (28)$$

Where $P_4^* = \int_0^{\infty} p_4^*(a_4)\, da_4$

Equation (28) yields $P^* = g^{-4}(b)$,

From (26),

$p_4^*(a) = p_4^*(0)e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}$

Substituting into (27), we need to solve

$$\left(1 - h_4(Q_4^*)\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}da_4\right)p_4^*(0) = 0$$

For $Q_4^*$

Fix $p_4^*(0)$ and $P_4^*$ Let

$$H_4(z_4) = 1 - h_4(z_4)\int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}da_{,4}$$

And assume that

$$[k_3]_4\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(g_4^{-4}(b_4))a_4 - [k_2]_4 a_4}\, da_4 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_4(z_4)] > 0.$$

Clearly,

$$\frac{dH_4}{dz_4} = h_4(z)(\,[f_2']_4(z_4)\,)\int_0^\infty a_4\beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - f_4(P_4^*)a_4 - f_2(z_4)a_4}da_{\square}$$

$-h'(z_4)\int_0^\infty \beta_4(a_4)\, e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}\quad da_4$

the assumptions on $\left[f_2^4\right]_4 and\ [h^4]^4 dH_4/dz > 0$. Hence, if

$$H_4(0) = 1 - \int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4} da_4 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_4(0) < 0$ follows if $g_4^-(b_4) < [f_4^{-4}]_4(r_4)$, where $r_4$ is the solution of (12).

Indeed, since

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 e^{-y\frac{a}{4}} da_4}$$

Is a decreasing function of $s_4$ when $P_4^* = [g^{-4}]_4(b_4) < [f_4^{-4}]_4(r_4)$, that is $[f_4]_4(P_4^*) < r$, $R\left(f_4(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_4(g^{-4}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4\left(g_4^{-4}(b)\right)a_4} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_4^*(a_4), Q_4^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_4^*(a_4), Q_4^*)$ is more complicated. We again consider perturbations $m_4(t_4 a_4) = p_4(t_4, a_4) - p_4^*(a_4)$ and $W_4(t) = Q_4(t) - Q^*$.Then

$$M_4(t) \equiv \int_0^\infty m_4(t, a_4) da_4 = P_4(t) - P_4^*,$$

$$[m_4]_4 + [m_a]_4 = [p_4]_4 + [p_a]_4 - [p^*]_4(a_4) \approx -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4$$

$$-[f_{4*}^4]_4 p_4^*(a_4)M - [f_{2*}^4]_4 p_4^*(\mathrm{a})W_4,$$

$$m_4(t,0)p_4(t,0) - p_4^*(0) \approx h_* \int_0^\infty \beta_4(a_4)m_4(t,a_4)\,da + \frac{[h_*^4]_4}{[h_*]_4}\, p_4^*(0)W_4$$

$$\frac{dW_4}{dt} = \frac{dQ_4}{dt} \approx [g_*^4]_4 Q_4^* M_{,4}$$

Where $[f_{4*}]_4$ denoting $[f_4(P^*)]_4$, $[f_{4*}]_4$ denoting $[f_2(P^*)]_4$, $[f_{4*}^{'}]_4$ denoting $[f^{'}{}_4(P^*)]_4 [f_{2*}^{'}]_4$ denoting $[f_2]_4$ $(Q_4^*,$ $h_{*}^{'}$ denoting $h_*^{'}$ denoting $h_4^{'}$ $(Q_4^*)$, and $[g_*^{'}]_4$

Denoting $[g^{'}]_4(P_4^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_4]_4 + [m_a]_4 = -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4 - [f^{'}{}_{4*}]_4 [p_4^*](a_4)M_4 - [f_{2*}^{'}]_4 p_4^*(a_4)W_4$$

$$m_4(t,0) = [h_*]_4 \int_0^\infty \beta_4(a_4)m_4(t,a_4)da_4 + \frac{[h_*^4]_4}{[h^*]_4}[p^*]_4(0)\, W_{,4} \qquad (30)$$

$$\frac{dW_4}{dt} = [g_*^4]_4 [Q_4^*]M_4$$

This is equation is important in the following cases: problem of using observed correlations to infer causal relations is relevant to (e) a wide variety of scientific disciplines. Yet given correlations between just two classical variables, (eb) it is impossible to determine whether they arose from (e) a causal influence of one on the other or a common cause influencing both, unless one can implement (eb) a randomized intervention. **Katja Ried, Megan Agnew,**

**Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. Resch** here consider the problem of causal inference for (e) quantum variables. Authors introduce causal tomography, which unifies and generalizes (e&eb) conventional quantum tomography schemes to provide (eb) a complete solution to the causal inference problem using (e) a quantum analogue of a randomized trial. Authors furthermore show that, in contrast to the classical case, observed quantum correlations alone can sometimes provide (eb) a solution. They implement a quantum-optical experiment that allows (eb) to control (e&eb) the causal relation between two optical modes, and two measurement schemes -- one with (e&eb) and one without (e) randomization -- that extract this relation from (e) the observed correlations. Results show that entanglement and (e&eb) coherence, known to be central to (e) quantum information processing, also provide (eb) a quantum advantage for (e) causal inference. Subjects: Quantum Physics (quant-ph); Learning (cs.LG); General Relativity and Quantum Cosmology (gr-qc); Machine Learning (stat.ML) Journal reference: Nat Phys 11, 414-420 (2015) DOI: 10.1038/nphys3266 Cite as: arXiv: 1406.5036 [quant-ph] (or arXiv: 1406.5036v1 [quant-ph] for this version) **Inferring causal structure: a quantum advantage Katja Ried, Megan Agnew, Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. ReschSTOP D. T. Son, A. O. Starinets** review recent progress in applying the AdS/CFT correspondence to (e&eb) finite-temperature field theory. In particular, they show how the hydrodynamic behavior of field theory is reflected in (e&eb) the low-momentum limit of (e) correlation functions computed through (e&eb) a real-time AdS/CFT prescription, which they formulate. Authors also show how the hydrodynamic modes in field theory correspond to (e&eb) the low-lying quasinormal modes of (e) the AdS black p-brane metric. They provide a proof of the universality of (e) the viscosity/entropy ratio within (eb) a class of theories with (e&eb) gravity duals and formulate (eb) a viscosity bound conjecture. Possible implications for real systems are mentioned. Subjects: High Energy Physics - Theory (hep-th) Journal reference: Ann.Rev.Nucl.Part.Sci.57:95-118,2007 DOI: 10.1146/annurev.nucl.57.090506.123120Report number: INT PUB 07-02Cite as: arXiv:0704.0240 [hep-th]

(or arXiv:0704.0240v2 [hep-th] for this version) An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society **Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP** Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv:

1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version **Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP Viscosity, Black Holes, and Quantum Field Theory D. T. Son, A. O. Starinets STOP Authors** study information retrieval from evaporating black holes, assuming that the internal dynamics of a black hole is unitary and rapidly mixing, and assuming that the retriever has unlimited control over the emitted Hawking radiation. If the evaporation of the black hole has already proceeded past the ``half-way" point, where (e) half of the initial entropy has been radiated away, then (eb) additional quantum information deposited in (eb) the black hole is revealed in (eb) the Hawking radiation very rapidly. Information deposited prior to the half-way point remains concealed until (e) the half-way point, and then emerges (eb) quickly. These conclusions hold because (e) typical local quantum circuits are efficient encoders for (e) quantum error-correcting codes that nearly achieve (eb) the capacity of the quantum erasure channel. Estimate of a black hole's information retention time, based on (e) speculative dynamical assumptions, is just barely compatible with (e&eb) the black hole complementarity hypothesis. **Black holes as mirrors: quantum information in random subsystems Patrick Hayden1 and John Preskill2 Published 26 September 2007 Journal of High Energy Physics, Volume 2007, JHEP09 (2007) STOP** By analyzing the key properties of black holes from the point of (e) view of quantum information, authors derive a model-independent picture of (e) black hole quantum computing. It has been noticed that this picture exhibits (eb) striking similarities with (e&eb) quantum critical condensates, allowing (eb) the use of a common language to describe (eb) quantum computing in both systems. Authors analyze such quantum computing by allowing coupling to (e&eb) external modes, under the condition that the external influence must be (=) soft-enough in order not to (e) offset the basic properties of the system. Authors derive model-independent bounds on some crucial time-scales, such as (=) the times of gate operation, decoherence, maximal entanglement and total scrambling. They show that for black hole type quantum computers all these time-scales are of (e) the order of the

black hole half-life time. Furthermore, authors construct explicitly a set of Hamiltonians that generates (eb) a universal set of quantum gates for (e) the black hole type computer. They find that the gates work at maximal energy efficiency. Furthermore, authors establish a fundamental bound on the complexity of quantum circuits encoded on (eb&e) these systems, and characterize the unitary operations that are (=) implementable. It becomes apparent that the computational power is very limited due to (e) the fact that the black hole life-time is of the same order of (e) the gate operation time. As a consequence, it is impossible to retrieve its information, within (eb) the life-time of a black hole, by externally coupling to (e&eb) the black hole qubits. However, authors show that, in principle, coupling to (e&eb) some of the internal degrees of freedom allows (eb) acquiring knowledge about the micro-state. Still, due to the trivial complexity of operations that can be performed, (eb) there is no time advantage over the collection of Hawking radiation and subsequent decoding. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Report number: LMU-ASC 03/16, MPP-2016-1 Cite as: arXiv:1605.01407 [hep-th] (or arXiv:1605.01407v1 [hep-th] for this version) **Universality of Black Hole Quantum Computing Gia Dvali, Cesar Gomez, Dieter Lust, Yasser Omar, and Benedikt Richter (Submitted on 4 May 2016) STOP** In this addendum to [arXiv: 1402.5674] two points are discussed. In the first additional evidence is provided (eb) for a dual connection between the geometric length of an Einstein-Rosen bridge and (e&eb) the computational complexity of the quantum state of the dual CFT's. The relation between growth of complexity and (e&eb) Page's “Extreme Cosmic Censorship” principle is also remarked on. The second point involves a gedanken experiment in which Alice measures (eb) a complete set of commuting observables at her end of (e) an Einstein-Rosen bridge is discussed. An apparent paradox is resolved by appealing to (e) the properties of GHZ tripartite entanglement. **Addendum to computational complexity and black hole horizons Leonard Susskind* Version of Record online: 18 JAN 2016 DOI: 10.1002/prop.201500093 STOP** Analyzing some well established facts, authors give a model-independent

parameterization of black hole quantum computing in terms of (e&*eb) a set of macro and micro quantities and their relations. These include the relations between the extraordinarily-small energy gap of (e) black hole qubits and important time-scales of (e) information-processing, such as, scrambling time and Page's time. Authors then show, confirming and extending previous results, that other systems of nature with (e&eb) identical quantum informatics features are (=) attractive Bose-Einstein systems at (eb) the critical point of quantum phase transition. Here authors establish a complete isomorphy between (e&eb) the quantum computational properties of these two systems. In particular, they show that the quantum hair of a critical condensate is (=) strikingly similar to (=) the quantum hair of a black hole. Irrespectively whether one takes the similarity between the two systems as a remarkable coincidence or as (=) a sign of a deeper underlying connection, the following is evident. Black holes are not unique in their way of (e) quantum information processing and they can manufacture black hole based quantum computers in labs by (e) taking advantage of quantum criticality. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Atomic Physics (physics. Atom-ph); Quantum Physics (quant-ph) Cite as: arXiv: 1601.01329 [hep-th] (or arXiv: 1601.01329v1 [hep-th] for this version) **Black Hole Based Quantum Computing in Labs and in the Sky Gia Dvali, Mischa Panchenko (Submitted on 6 Jan 2016) STOP Gia Dvali, Alexander Gußmann** show that the existence of black holes with (e&eb) classical skyrmion hair invalidates (e) standard proofs that global charges, such as (=) the baryon number, cannot be conserved by (e) a black hole. By carefully analyzing the standard arguments based on (e) a Gedankenexperiment in which a black hole is seemingly-unable to (e) return the baryon number that it swallowed, authors identify inconsistencies in this reasoning, which does not (e) take into the account neither the existence of skyrmion black holes nor the baryon/skyrmion correspondence. Authors then perform a refined Gedankenexperiment by incorporating (e) the new knowledge and show that (eb) no contradiction with conservation of baryon number takes place at (eb) any stage of black hole evolution. Analysis also indicates (eb) no conflict

between semi-classical black holes and (e&eb) the existence of baryonic gauge interaction arbitrarily-weaker than gravity. Next, authors study classical cross sections of a minimally-coupled massless probe scalar field scattered by (e) _ a skyrmion black hole. Authors investigate how the skyrmion hair manifests itself by (e) comparing this cross section with (e&eb) the analogous cross section caused by (e) a Schwarzschild black hole which has (e) the same ADM mass as (=) the skyrmion black hole. Here authors find an order-one difference in the positions of (e) the characteristic peaks in (eb) the cross sections. The peaks are shifted to smaller scattering angles when (e) the skyrmion hair is present. This comes from the fact that the skyrmion hair changes (e&eb) the near horizon geometry of the black hole when compared to (e&eb) a Schwarzschild black hole with same ADM mass. Authors keep the study of this second aspect general so that (e) the qualitative results obtained can also be applied to(e&eb) black holes with classical hair of different kind. Subjects: High Energy Physics - Theory (hep-th); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU-ASC 18/16 Cite as: arXiv:1605.00543 [hep-th] (or arXiv:1605.00543v2 [hep-th] for this version) A preliminary discussion is given of (e) the prospects that gravitational-wave observations of (e) binary inspiral of black holes could reveal or constrain (e&eb) quantum modifications to (e) black hole dynamics, such as are required to (e) preserve postulates of quantum mechanics. Different proposals for such modifications are characterized by (e) different scales, and the size of these scales relative to (e&eb) those probed by observation of (e) inspiral signals is important in determining (eb) the feasibility of finding experimental signatures. Certain scenarios with strong quantum modifications in (eb) a region extending well outside the horizon are expected to modify (e&eb) classical evolution, and distort (e&eb) the near-peak gravitational wave signal, suggesting (eb) a search for departures from (e) waveforms predicted by general relativity. The near agreement of the GW150914 signals with (e&eb) such waveforms is discussed, and indicates (eb) constraints on some such scenarios. Important strategies for more precise future tests are 1) to develop more precise

predictions from (e) scenarios proposing quantum modifications, and 2) searching for observed deviations from (e) numerical relativity predictions via (e&eb) analysis of gravity wave data, particularly focussing on (e&eb) the signal region corresponding to plunge and merger. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1602.03622 [gr-qc] (or arXiv: 1602.03622v2 [gr-qc] for this version) **Gravitational wave tests of quantum modifications to black hole structure -- with post-GW150914 update Steven B. Giddings STOP** Direct LHC bounds on (e&eb) colored SUSY particles now corner naturalness more than the measured value of the Higgs mass does. Bounds on the gluino are of (e) particular importance, since it radiatively "sucks" up (e) the stop and Higgs-up soft masses. As a result, even models that easily accommodate a 125 GeV Higgs are almost as tuned as (=) the simplest version of SUSY, the MSSM: at best at the percent level. In this paper, authors further examine how current LHC results constrain (e) naturalness in three classes of models that may relax LHC bounds on (e&eb) sparticles: split families (e&eb), baryonic RPV (e&eb), and Dirac gauginos. In models of split families and bRPV, the bounds on the gluino are (=) only slightly reduced, resulting in (eb) a few percent tuning. In particular, having a natural spectrum in bRPV models typically implies (eb) that tops, Ws, and Zs are easily produced (eb) in the cascade decays of (e) squarks and gluinos. The resulting leptons and missing energy push (eb) the gluino mass limit above 1 TeV. Even when the gluino has a Dirac mass and does not contribute to the stop mass at one loop, tuning reappears in calculable models because there is no symmetry imposing the supersoft limit. Authors conclude that, even if sparticles are found at LHC-14, naturalness will not emerge triumphant. Subjects: High Energy Physics - Phenomenology (hep-ph)DOI:10.1007/JHEP03(2014)022 Cite as:rXiv:1309.3568 [hep-ph] (or arXiv:1309.3568v2 [hep-ph] for this version) **The Last Vestiges of Naturalness Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang, Ken Van Tilburg, Giovanni Villadoro STOP** The lack of evidence for new physics beyond (e) the standard model at the LHC points to (e) a paucity of new particles near the weak scale. This suggests that the weak scale is tuned and that supersymmetry, if present at

all, is realized (eb) at higher energies. The measured Higgs mass constrains (e) the scalar sparticles to be below 10^5 TeV, while (e) gauge coupling unification favors (eb) Higgsinos below 100 TeV. Nevertheless, in many models gaugino masses are (=) suppressed and remain within reach of the LHC. Tuning the weak scale and (e&eb) the renormalization group evolution of (e) the scalar masses constrain (e) Split model building. Due to the small gaugino masses, either the squarks or the up-higgs often run tachyonic; in the latter case, successful electroweak breaking requires heavy higgsinos near the scalar sparticles. We discuss the consequences of tuning the weak scale and the phenomenology of several models of Split supersymmetry including anomaly mediation, U(1)_(B-L) mediation, and Split gauge mediation. Subjects: High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Experiment (hep-ex) DOI: 10.1007/JHEP02 (2013)126 Cite as: arXiv: 1210.0555 [hep-ph] (or arXiv: 1210.0555v2 [hep-ph] for this version) **Mini-Split Asimina Arvanitaki, Nathaniel Craig, Savas Dimopoulos, Giovanni Villadoro STOP** Light states associated with (e&eb) the hierarchy problem affect (e&eb) the Higgs LHC production and decays. Authors illustrate this within the MSSM and two simple extensions applying (e&eb) the latest bounds from (e) LHC Higgs searches. Large deviations in the Higgs properties are expected in (eb) a natural SUSY spectrum. The discovery of a non-Standard-Model Higgs may signal (eb) the presence of light stops accessible at (eb) the LHC. Conversely, the more the Higgs is Standard-Model-like, (eb) the more tuned the theory becomes. Taking the ratio of different Higgs decay channels at the LHC cancels (er&eb) the leading (eb) QCD uncertainties and potentially improves (eb) the accuracy in Higgs coupling measurements to (e) the percent level. This may lead to (eb) the possibility of doing precision Higgs physics at (eb) the LHC. Finally, authors entertain the possibility that the ATLAS excess around (e&eb) 125 GeV persists with (e&eb) a Higgs production cross-section that is enhanced compared to (e&eb) the SM. This increase can only be accommodated in extensions of (e) the MSSM and it may suggest (eb) that stops lie below 400 GeV, likely within (eb) reach of next year's LHC run. Subjects: High Energy Physics - Phenomenology (hep-ph) DOI: 10.1007/JHEP02 (2012)144 Cite

as: arXiv: 1112.4835 [hep-ph] (or arXiv: 1112.4835v2 [hep-ph] for this version) **A Non Standard Model Higgs at the LHC as a Sign of Naturalness Asimina Arvanitaki, Giovanni Villadoro STOP** It has recently been suggested that the presence of a plenitude of light axions, an Axiverse, is (=) evidence for the extra dimensions of string theory. Authors discuss the observational consequences of (e) these axions on astrophysical black holes through (e&eb) the Penrose superradiance process. When an axion Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e&eb) the black hole "nucleus" forming (eb) a gravitational atom in the sky. The occupation number of superradiant atomic levels fed by (e) the energy and angular momentum of the black hole grows exponentially. The black hole spins down and an axion Bose-Einstein condensate cloud forms (eb) around it. When the attractive axion self-interactions become (=) stronger than the gravitational binding energy, (eb) the axion cloud collapses, a phenomenon known in condensed matter physics as "Bosenova". The existence of axions is first diagnosed by (e) gaps in the mass vs (e&eb) spin plot of astrophysical black holes. For young black holes the allowed values of (e) spin are quantized, giving rise to (eb) "Regge trajectories" inside the gap region. The axion cloud can also be observed directly either through (e&eb) precision mapping of the near horizon geometry or through (e&eb) gravitational waves coming from (e) the Bosenova explosion, as well as axion transitions and (e&eb) annihilations in (eb) the gravitational atom. Estimates suggest (eb) that these signals are detectable in (eb) upcoming experiments, such as Advanced LIGO, AGIS, and LISA. Current black hole spin measurements imply (eb) an upper bound on (e&eb) the QCD axion decay constant of (e) 2 x 10^17 GeV, while (e) Advanced LIGO can detect (e&eb) signals from (e) a QCD axion cloud with (e&eb) a decay constant as low as (=) the GUT scale. Authors finally discuss the possibility of observing the gamma-rays associated with the Bosenova explosion and, perhaps, the radio waves from axion-to-photon conversion for the QCD axion. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Journal

reference:Phys.Rev.D83:044026, 2011 DOI: 10.1103/PhysRevD.83.044026 Cite as: arXiv: 1004.3558 [hep-th] (or arXiv: 1004.3558v2 [hep-th] for this version) **Exploring the String Axiverse with Precision Black Hole Physics Asimina Arvanitaki, Sergei Dubovsky (Submitted on 20 Apr 2010 (v1), last revised 16 Aug 2011 (this version, v2)) STOP Skyrmion Black Hole Hair: Conservation of Baryon Number by Black Holes and Observable Manifestations Gia Dvali, Alexander Gußmann STOP Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust** discuss BMS supertranslations both at (eb) null-infinity and on the horizon for the case of the Schwarzschild black hole. Authors show that both kinds of supertranslations lead to (eb) infinitely many gapless physical excitations. On this basis authors construct a quotient algebra using (e) suited superpositions of (e) both kinds of transformations which cannot be compensated by (e) an ordinary BMS-supertranslation and therefore are (=) intrinsically due to (e) the presence of an event horizon. Authors show that these quotient transformations are (=) physical and generate (eb) gapless excitations on the horizon that can account for (e) the gravitational hair as well as for (e) the black hole entropy. They identify the physics of these modes as associated with (e&eb) Bogolioubov-Goldstone modes due to quantum criticality. Classically the number of these gapless modes is (=) infinite. However, authors show that due to (e) quantum criticality the actual amount of information-carriers becomes (=) finite and consistent with Bekenstein entropy. Although they only consider the case of Schwarzschild geometry, the arguments are extendable to(e&eb) arbitrary space-times containing event horizons. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU--ASC 05/16, MPP--2016—3 Cite as: arXiv: 1601.03725 [hep-th] (or arXiv: 1601.03725v1 [hep-th] for this version) **Gravitational Black Hole Hair from Event Horizon Supertranslations Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust STOP Gia Dvali, Cesar Gomez, Nico Wintergerst** suggest that holography can be formulated in terms of (e&eb) the information capacity of (e) the St\"uckelberg degrees of freedom that maintain

gauge invariance of (e) the theory in the presence of (e) an information boundary. These St\"uckelbergs act as (e&eb) qubits that account for a certain fraction of (e) quantum information. Their information capacity is measured by (e) the ratio of the inverse St\"uckelberg energy gap to (e) the size of the system. Systems with (e&eb) the smallest gap are (=) maximally holographic. For massless gauge systems this information measure is (=) universally equal to (=) the inverse coupling evaluated at the systems' length scale. In this language it becomes very transparent why the St\"uckelberg information capacity of black holes saturates (e full) the Bekenstein bound and accounts for (e) the entire information of the system. The physical reason is that the strength of quantum interaction is bounded from below by the gravitational coupling, which scales as (=) area. Observing the striking similarity between the scalings of the energy gap of the boundary St\"uckelberg modes and (e&eb): (AC=DC=0) the Bogoliubov modes of critical many-body systems, authors establish a connection between holography and (e&eb) quantum criticality through (e&eb) the correspondence between these modes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1511.03525 [hep-th] (or arXiv: 1511.03525v1 [hep-th] for this version) **Stückelberg Formulation of Holography Gia Dvali, Cesar Gomez, Nico Wintergerst STOP B**lack hole entropy, denoted by N, in (semi)classical limit is (=) infinite. This scaling reveals (eb) a very important information about the qubit degrees of (e) freedom that carry black hole entropy. Namely, the multiplicity of qubits scales as N, whereas (e) their energy gap and their coupling as 1/N. Such a behavior is indeed exhibited by (e) Bogoliubov-Goldstone degrees of freedom of (e) a quantum-critical state of N soft gravitons (a condensate or a coherent state) describing (eb) the black hole quantum portrait. They can be viewed as (=) the Goldstone modes of (e) a broken symmetry acting on (e&eb) the graviton condensate. In this picture Minkowski space naturally emerges as (=) a coherent state of infinite-N gravitons of (e) infinite wavelength and it carries (e&eb) an infinite entropy. In this paper authors **Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015)** ask what is the geometric meaning (if any) of (e) the classical limit of this symmetry. Authors argue that the infinite-

N limit of Bogoliubov-Goldstone modes of critical graviton condensate is described by (e) recently-discussed classical BMS super-translations broken by the black hole geometry. However, the full black hole information can only be recovered for (e) finite N, since the recovery time becomes (=) infinite in classical limit in which N is (=) infinite. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) DOI: 10.1016/j.physletb.2015.11.073 Report number: MPP-2015-215, LMU-ASC 59/15 Cite as: arXiv: 1509.02114 [hep-th] (or arXiv: 1509.02114v1 [hep-th] for this version) **Classical Limit of Black Hole Quantum N-Portrait and BMS Symmetry Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015) STOP** Recent ideas about understanding physics of black hole information-processing in terms of (e&eb) quantum criticality allow us to (eb) implement black hole mechanisms of quantum computing within (eb) critical Bose-Einstein systems. The generic feature, uncovered both by analytic and numeric studies, is (=) the emergence at the critical point of (e) gapless weakly-interacting modes, which act as (=) qubits for information-storage at a very low energy cost. These modes can be effectively described in terms of (e&eb) either Bogoliubov or Goldstone degrees of freedom. The ground-state at the critical point is (=) maximally entangled and far from being classical. We confirm this near-critical behavior by a new analytic method. **Gia Dvali, Mischa Panchenko** compute growth of entanglement and show its consistency with (e&eb) black hole type behavior. On the other hand, in the over-critical regime the system develops (eb) a Lyapunov exponent and scrambles quantum information very fast. By, manipulating (e&eb) the system parameters externally, authors can put it in and out of various regimes and in this way control (e&eb) the sequence of information storage and processing. By using gapless Bogoliubov modes as (=) control qubits they design some simple logic gates. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Cite as: arXiv: 1507.08952 [hep-th] (or arXiv: 1507.08952v1 [hep-th] for this version) **Black Hole Type Quantum**

**Computing in Critical Bose-Einstein Systems Gia Dvali, Mischa Panchenko (Submitted on 31 Jul 2015) STOP Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst** establish a fundamental connection between quantum criticality of a many-body system, such as Bose-Einstein condensates, and (e&eb) its capacity of information-storage and processing. For deriving the effective theory of modes in the vicinity of the quantum critical point they develop a new method by mapping a Bose-Einstein condensate of N-particles onto (e&eb) a sigma model with a continuous global (pseudo)symmetry that mixes (e&eb) bosons of different momenta. The Bogolyubov modes of the condensate are mapped onto (e&eb) the Goldstone modes of the sigma model, which become (=) gapless at the critical point. These gapless Goldstone modes are (=) the quantum carriers of information and entropy. Analyzing their effective theory, authors observe the information-processing properties strikingly similar to (=) the ones predicted by the black hole portrait. The energy cost per qubit of information-storage vanishes in (eb) the large-N limit and the total information-storage capacity increases with (e&eb) N either exponentially or as a power law. The longevity of information-storage also increases with (e&eb) N, whereas (e) the scrambling time in (eb) the over-critical regime is controlled by (e&eb) the Lyapunov exponent and scales logarithmically with (e&eb) N. This connection reveals (eb) that the origin of black hole information storage lies in (eb) the quantum criticality of the graviton Bose-gas, and that much simpler systems that can be manufactured in table-top experiments can exhibit (eb) very similar information-processing dynamics. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 92, 125002 (2015) DOI: 10.1103/PhysRevD.92.125002 Cite as: arXiv:1507.02948 [hep-th] (or arXiv:1507.02948v1 [hep-th] for this version) **Nambu-Goldstone Effective Theory of Information at Quantum Criticality Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst STOP Gia Dvali, Cesar Gomez** show that a positive cosmological constant is (=) incompatible with the quantum-corpuscular resolution of de Sitter metric in form of (e) a coherent state. The reason is very general and is due to (e) the

quantum self-destruction of (e) the coherent state because of (e) the scattering of constituent graviton quanta. This process creates an irreversible quantum clock, which precludes (e) eternal de Sitter. It also eliminates (e) the possibility of Boltzmann brains and Poincare recurrences. This effect is expected to be (=) part of any microscopic theory that takes into account (e) the quantum corpuscular structure of the cosmological background. This observation puts the cosmological constant problem in a very different light, promoting it, from a naturalness problem, into (e&eb) a question of quantum consistency. We are learning that quantum gravity cannot tolerate (e) exceedingly-classical sources. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1412.8077 [hep-th]

(or arXiv:1412.8077v1 [hep-th] for this version) **Quantum Exclusion of Positive Cosmological Constant? Gia Dvali, Cesar Gomez (Submitted on 27 Dec 2014) STOP G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger** establish a connection between the ultra-Planckian scattering amplitudes in field and string theory and (e&eb) unitarization by black hole formation in these scattering processes. **Using as a guideline an explicit microscopic theory in which the black hole represents (eb) a bound-state of many soft gravitons at (eb) the quantum critical point, they were able to identify and compute (e&eb) a set of perturbative amplitudes relevant for (e) black hole formation. These are the tree-level N-graviton scattering S-matrix elements in (eb) a kinematical regime (called classicalization limit) where (e) the two incoming ultra-Planckian gravitons produce (eb) a large number N of soft gravitons. Authors compute these amplitudes by using (e) the Kawai-Lewellen-Tye relations, as well as scattering equations and (e&eb) string theory techniques. Authors discover that this limit reveals (eb) the key features of the microscopic corpuscular black hole N-portrait. In particular, the perturbative suppression factor of a N-graviton final state, derived from (e) the amplitude, matches the non-perturbative black hole entropy when (e) N reaches the quantum criticality value, whereas final states with different value of N are (=) either suppressed or excluded by non-perturbative**

**corpuscular physics. Thus they identify the microscopic reason behind the black hole dominance over other final states including (e) non-black hole classical object. In the parameterization of the classicalization limit the scattering equations can be solved exactly allowing (eb) us to obtain closed expressions for the high-energy limit of the open and closed superstring tree-level scattering amplitudes for a generic number N of external legs. Authors demonstrate matching and complementarity between the string theory and (e&eb) field theory** in different large-s and large-N regimes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1016/j.nuclphysb.2015.02.004 Report number: MPP-2014-320, LMU-ASC 52/14 Cite as: arXiv:1409.7405 [hep-th] (or arXiv:1409.7405v2 [hep-th] for this version) **Black Hole Formation and Classicalization in Ultra-Planckian 2 -> N Scattering G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger STOP** Gravitational backgrounds, such as black holes, AdS, de Sitter and inflationary universes, should be viewed as (=) composite of N soft constituent gravitons. It then follows that such systems are close to (e) quantum criticality of graviton Bose-gas to Bose-liquid transition. Generic properties of the ordinary metric description, including (e) geodesic motion or particle-creation in the background metric, emerge as (=) the large-N limit of quantum scattering of (e) constituent longitudinal gravitons. **Gia Dvali, Cesar Gomez** show that this picture correctly accounts for (e) physics of large and small black holes in AdS, as well as reproduces (eb) well-known inflationary predictions for cosmological parameters. However, it anticipates new effects not captured by (e) the standard semi-classical treatment. In particular, **Gia Dvali, Cesar Gomez** predict observable corrections that are sensitive to (e&eb) the inflationary history way beyond last 60 e-foldings. They derive an absolute upper bound on the number of e-foldings, beyond (e) which neither de Sitter nor inflationary Universe can be approximated by (e) a semi-classical metric. However, they could in principle persist in a new type of (e) it quantum eternity state. Authors discuss implications of this phenomenon for (e&eb) the cosmological constant problem. Subjects: High Energy Physics - Theory (hep-th); Cosmology and

Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Quantum Physics (quant-ph) DOI: 10.1088/1475-7516/2014/01/023 Cite as: arXiv: 1312.4795 [hep-th] (or arXiv:1312.4795v1 [hep-th] for this version) **Quantum Compositeness of Gravity: Black Holes, AdS and Inflation Gia Dvali, Cesar Gomez STOP** As quantum information science approaches the goal of constructing quantum computers, understanding loss of information through (e&eb) decoherence becomes (=) increasingly important. The information about a system that can be obtained from (e) its environment can facilitate quantum control and error correction. Moreover, observers gain (eb) most of their information indirectly, by (e) monitoring (primarily photon) environments of (e) the "objects of interest." Exactly how this information is inscribed in the environment is essential for (e) the emergence of "the classical" from (e) the quantum substrate. In this paper, **Robin Blume-Kohout, Wojciech H. Zurek** examine how many-qubit (or many-spin) environments can store (e) information about a single system. The information lost to (e) the environment can be stored redundantly, or it can be encoded in (eb) entangled modes of the environment. Authors go on to show that randomly chosen states of the environment almost always encode (e) the information so that (e) an observer must capture (e) a majority of the environment to deduce (eb) the system's state. Conversely, in the states produced by (e) a typical decoherence process, information about a particular observable of the system is stored redundantly. This selective proliferation of "the fittest information" (known as Quantum Darwinism) plays a key role in (eb) choosing the preferred, effectively classical observables of (e) macroscopic systems. The developing appreciation that the environment functions not just as (=) a garbage dump, but as a communication channel, is extending understanding of the environment's role in (e&eb) the quantum-classical transition beyond (e) the traditional paradigm of decoherence. Keywords Decoherence quantum Darwinism quantum information quantum-classical transition **Foundations of Physics November 2005, Volume 35, Issue 11, pp 1857-1876 A Simple Example of "Quantum Darwinism": Redundant Information Storage in Many-Spin Environments Robin Blume-Kohout, Wojciech H. Zurek STOP Harold Ollivier, David**

**Poulin, and Wojciech H. Zurek** study the emergence of objective properties in open quantum systems. In our analysis, the environment is promoted from a passive role of a reservoir selectively destroying quantum coherence to an active role of amplifier selectively **proliferating information about the system. We show that only preferred pointer states of the system can leave a redundant and therefore easily detectable imprint on the environment. Observers who—as is almost always the case—discover the state of the system indirectly (by probing a fraction of its environment) will find out only about the corresponding pointer observable. Many observers can act in this fashion independently and without perturbing the system. They will agree about its state. In this operational sense, preferred pointer states exist objectively. Received 24 July 2003 DOI:http://dx.doi.org/10.1103/PhysRevLett.93.220401 Objective Properties from Subjective Quantum States: Environment as a Witness Harold Ollivier, David Poulin, and Wojciech H. Zurek Phys. Rev. Lett. 93, 220401 – Published 22 November 2004 STOP** Quantum Darwinism recognizes that we—the observers—acquire our information about the "systems of interest" indirectly from their imprints on the environment. **Michael Zwolak, H. T. Quan, and Wojciech H. Zurek** show that information about a system can be acquired from a mixed-state, or hazy, environment, but the storage capacity of an environment fragment is suppressed by its initial entropy. In the case of good decoherence, the mutual information between the system and the fragment is given solely by the fragment's entropy increase. For fairly mixed environments, this means a reduction by a factor 1−h, where h is the haziness of the environment, i.e., the initial entropy of an environment qubit. Thus, even such hazy environments eventually reveal the state of the system, although now the intercepted environment fragment must be larger by ~ (1−h)−1 to gain the same information about the system. DOI:http://dx.doi.org/10.1103/PhysRevLett.103.110402 **Quantum Darwinism in a Mixed Environment Michael Zwolak, H. T. Quan, and Wojciech H. Zurek Phys. Rev. Lett. 103, 110402 – Published 8 September 2009** We study quantum Darwinism—the redundant recording of information about the preferred states of a decohering system by its environment—for an object illuminated by a

blackbody. In the cases of point-source and isotropic illumination, we calculate the quantum mutual information between the object and its photon environment. We demonstrate that this realistic model exhibits fast and extensive proliferation of information about the object into the environment and results in redundancies orders of magnitude larger than the exactly soluble models considered to date. DOI:http://dx.doi.org/10.1103/PhysRevLett.105.020404 **Quantum Darwinism in an Everyday Environment: Huge Redundancy in Scattered Photons C. Jess Riedel and Wojciech H. Zurek Phys. Rev. Lett. 105, 020404 – Published 9 July 2010 STOP** Measurements transfer information about a system to the apparatus and then, further on, to observers and (often inadvertently) to the environment. I show that even imperfect copying essential in such situations restricts possible unperturbed outcomes to an orthogonal subset of all possible states of the system, thus breaking the unitary symmetry of its Hilbert space implied by the quantum superposition principle. Preferred outcome states emerge as a result. They provide a framework for "wave-packet collapse," designating terminal points of quantum jumps and defining the measured observable by specifying its eigenstates. In quantum Darwinism, they are the progenitors of multiple copies spread throughout the environment—the fittest quantum states that not only survive decoherence, but subvert the environment into carrying information about them—into becoming a witness. DOI:http://dx.doi.org/10.1103/PhysRevA.76.052110 **Quantum origin of quantum jumps: Breaking of unitary symmetry induced by information transfer in the transition from quantum to classical Wojciech Hubert Zurek Phys. Rev. A 76, 052110 – Published 16 November 2007 STOP** The Born rule, a foundational axiom used to deduce probabilities of events from **wavefunctions, is indispensable in the everyday practice of quantum physics. It is also key in the quest to reconcile the ostensibly inconsistent laws of the quantum and classical realms, as it confers physical significance to reduced density matrices, the essential tools of decoherence theory. Following Bohr's Copenhagen interpretation, textbooks postulate the Born rule outright. However, recent attempts to derive it from (e) other quantum principles have been successful, holding promise for (e)**

**simplifying and clarifying the quantum foundational bedrock. A major family of derivations is based on envariance, a recently discovered symmetry of (e) entangled quantum states. Here, authors identify and experimentally test three premises central to (e) these envariance-based derivations, thus demonstrating, (eb) in the microworld, the symmetries from (e) which the Born rule is derived. Further, they demonstrate envariance in a purely local quantum system, showing (eb) its independence from relativistic causality.Subjects: Quantum Physics (quant-ph)Cite as:arXiv:1604.01471 [quant-ph] (or arXiv:1604.01471v1 [quant-ph] for this version) Quantum probabilities from quantum entanglement: Experimentally unpacking the Born rule Jérémie Harris, Frédéric Bouchard, Enrico Santamato, Wojciech H. Zurek, Robert W. Boyd, Ebrahim Karimi**

Integrating $[m_4]_4(t, a_4)$ with respect to a from 0 to ∞, we have

$$\frac{dM_4(t)}{dt} = \int_0^\infty [mt]_4(t, a_4)da_4$$

$$= -\int_0^\infty [m_a]_4\,(t, a_4)da_4 - \int_0^\infty \mu(a_4)m(t, a_4)da_4$$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f^{'}{}_{4*}]_4 \int_0^\infty p_4^*(a_4)da_4\right)M_4 - [f_{2*}^4]_4 \int_0^\infty p_4^*(a_4)da_4 W_4$$

$= m_4(t, 0) - \int_0^\infty \mu_4(a_4)\, m_4(t, a_4)da_4$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}^{i}]_4 P_4^*\right)M_4 - [f_{2*}^{i}]_4 P_4^* W_{\cdot 4}$$

From (30), we arrive at the equations

$$\frac{dW_4(t)}{dt} = \int_0^{\infty} ([h_*]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4$$

$$-\left([f_{4*}] + [f_{2*}]_4 + \left[f_{4*}^{'}\right]_4 P_4^*\right) M_4 + \left(\frac{\left[h_*^{'}\right]_4}{[h_*]^4}\, P_4^*(0) - f_{2*}^{'} P_4^*\right) W_4,$$

(34)

$$\frac{dW_4(t)}{dt}\left[g_{*}^{'}\right]_4 Q_4^*\, M_{.4}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{r*}]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4 = ([h_*]_4\beta_4(\theta_4) - \mu_4(\theta_4))M_{,4}$$

Where $0< \theta_4 = \theta_4(t) < \infty$.

And if $a_4^2\,[a_4^2]_4<4\,[a_2]_4<1\, a_2 e$,

$\mathrm{Re}\,\lambda_4\left[\lambda_4\right]_4 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_4 = -\frac{4}{2}a_4[a2\ ]_4 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_*\,[h_*]_4\beta_4(a_4) - \mu_4(a_4)\right|$. Then$|n_4(t)|$ ⊠⊿∀t ⊠. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_4^2 <4\, a_2 e$ and $[a_4] > 2\Delta$,or $[a_4^2]$ ⊠[$a_2$]e and $(a_4 - \sqrt{(a_4^2)_4 - 4(a_2)_4 e}) > 2\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+4}\right)$ and $\beta_4(a_4) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4) + d\tau_4 da} = \frac{4}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_4 + [\mu_0]$ 0.5. Assume$[f_4](x_4)= K_4 x_4$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_4(x_4)= [0x_4]/(1 + x_4)$; the functional responses $[[f_2](x) = [0]\, x_4/(1 + x_4)$and $h_4(x_4)=(1+\omega_4 x_4)/(1 + x_4)$(See Holling [7] and Rashevsky [13])

Equation (44),now has the form of

$$\frac{m_4 + [\mu_0] + [K][P_0])}{[m_4]_4 + (\mu_0)_4 + [K][P_0])\frac{3}{4}\,[K]_4[P_0]\frac{3}{4}} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(4 + P_0) < b_4$ i e $[P_0] < b/(\delta_4 - b_4)$then the equilibrium$(P_0)$ $[P_0]_4(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_4$

$[f_{4*}] = K_4\, P^*\, [f_{2*}] = \frac{\sigma Q_4^*}{4+Q_4^*} \quad [h_*] = \frac{4+\omega_4 Q_4^*}{4+Q_4^*}$

$[f_{4*}] = K_4\, [f_{2*}^4] = \frac{\sigma_4}{(4+Q_4^*)^2}\, [g_*^4] = \frac{\sigma_4}{(4+P_4^*)^2} \quad [h_*^4] = \frac{\omega_4 - 4}{(4+Q_4^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0$, K, b, $\delta$, and $\omega$ need to satisfy

$$\frac{\omega_4}{(m + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma)^2} + \frac{2\omega_4}{([m_4]_4 + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma_4)^3} < 1$$

And

$$\frac{1}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^2}+\frac{2\omega_4}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4}+(m_4+[\mu_0]_4+\frac{K_4b_4}{\delta_4-b_4}+\sigma_4)\,Q_4^*)^3$$

$$=m_4+[\mu_0]_4+3+\frac{K_4b_4}{\delta_4-b_4}(4+Q^*)(1+\omega_4Q_4^*$$

Now

$$\left[a_4^2\right]-1\left[a_2\right]e=\left(\frac{2K_4b_4}{\delta_4-b_4}\right)^2+\frac{1K_4b_4\sigma_4\,Q_4^*}{(\delta_4-b_4)(1+Q_4^*)}+\frac{\sigma_4^2Q_4^{*2}}{(1+Q_4^*)^2}$$

$$\frac{1\,b_4\sigma_4\,(\delta_4-b)Q^*}{\delta_4(1+Q_4^*)^2}-\frac{1\,(\delta_4-b_4)^2(1-\omega_4)Q_4^*}{\delta_4(1+\omega_4Q_4^*)(1+Q_4^*)}P^*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_4\,ae_4^{mu}\;[\mu_0]_4-\frac{a_4}{a_4+4}$$

Hence if $[a_4^2>4(a_2)_4e$ and

$$[a_4] = \frac{2K_4 b_4}{\delta_4 - b_4} + \frac{\sigma_4 \, Q_4^*}{1 + Q_4^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_4]_4 = [f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}] P_4^*$$

$$[a_2]_4 = \left[f_{2*}^4\right] P_4^* - \frac{h_*^4}{h_*} P^*(0)$$

$$e = [g_*^4]_4 \, Q_4^*$$

$$\eta_4(t) = \left[h_*\right]_4 \beta_4 \left(\theta_4(t)\right) - \mu_4(\theta(t)),$$

The system(34)can be written as the following system:

$$[M(t)]\left[a_4\right] M_4 - [a_2]_4 W_4 - [n(t)] M_4$$

$$w_4(t) = eM_4$$

Write

$$A = \begin{pmatrix} -[a_4]_4 - [a_2]_4 & \\ e & 0 \end{pmatrix}$$

$$B = (t) = \begin{pmatrix} [\eta(t)]_4 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_4 \\ W_4 \end{pmatrix}$$

The solution of (32)is

$$X_4(t) = e^{A,t}[X]_4(0) + \int_0 e^{At}$$

**Some of the other conceptual and thematic and discursive factorial representationalities to which the model is applicable are:every one-dimensional quantum mechanical Hamiltonian H1 can have (e) a partner H2 such that H1 and H2 taken together may be viewed as (=) the components of a supersymmetric Hamiltonian. The term 'supersymmetric Hamiltonian' is taken to mean (eb) a Hamiltonian defined in terms of (e&eb)charges that obey (eb)the same algebra as that of (e) the generators of supersymmetry in field theory. The consequences of this symmetry for (er) the spectra of H1 and H2 are explored. It is shown how the supersymmetric pairing may be utilised to eliminate (e) the ground state of H1, or add a state below the ground state of H1 or maintain the spectrum of H1. It is also explicitly demonstrated that the supersymmetric pairing may be used to generate (eb) a class of anharmonic potentials with (e&eb) exactly specified spectra. The complete spectrum of an anharmonic potential so generated consists of (e) all the eigenstates of the simple harmonic oscillator and, in addition, a ground state at a specified energy E which lies arbitrarily below (eb) the E=1/2 ground state of the harmonic oscillator. Supersymmetric quantum mechanics of one-dimensional systems C V Sukumar Journal of Physics A: Mathematical and General, Volume 18, Number 15 STOP Authors examine the properties and consequences of (e&eb) pseudo-supersymmetry for (e) quantum systems admitting (e) a pseudo-Hermitian Hamiltonian. Authors explore the Witten index of (e) pseudo-**

**supersymmetry and show that every pair of diagonalizable (not necessarily Hermitian) Hamiltonians with (e&eb) discrete spectra and (e&eb) real or complex-conjugate pairs of (e) eigenvalues are (=) isospectral and have (e) identical degeneracy structure except (e) perhaps for the zero eigenvalue if and only if (e) they are pseudo-supersymmetric partners. This implies (eb) that pseudo-supersymmetry is (=) the basic framework for (e) generating non-Hermitian PT-symmetric and (e&eb) non-PT-symmetric Hamiltonians with (e&eb) a real spectrum via (e&eb) a Darboux transformation, and shows (eb) that every diagonalizable Hamiltonian H with(e&eb) a discrete spectrum and (e&eb) real or complex-conjugate pairs of (e) eigenvalues may be factored as (=) H=L♯L where L is (=) a linear operator with pseudo-adjoint L♯. In particular, this factorization applies to (e&eb) PT-symmetric and Hermitian Hamiltonians. The non-degenerate two-level systems provide (eb) a class of Hamiltonians that are pseudo-Hermitian. Authors demonstrate the implications of general results for (e&eb) this class in some detail. Copyright © 2002 Elsevier Science B.V. All rights reserved. Nuclear Physics B Volume 640, Issue 3, 30 September 2002, Pages 419–434Pseudo-supersymmetric quantum mechanics and isospectral pseudo-Hermitian Hamiltonians Ali Mostafazadeh STOP Michele Caselle, Alessandro Nada, Marco Panero present a high-precision lattice calculation of (e) the equation of state in the confining phase of (e) SU(2) Yang-Mills theory. Authors show that the results are described very well by (e) a gas of massive, non-interacting glueballs, provided (e) one assumes an exponentially growing Hagedorn spectrum. The latter can be derived within (eb) an effective bosonic closed-string model, leading to (eb) a parameter-free theoretical prediction, which is in perfect agreement with (=) proposed lattice results. Furthermore, when applied to (e&eb) SU (3) Yang-Mills theory, this effective model accurately describes (eb) the lattice results reported by Bors\'anyi et al. in JHEP 07 (2012) 056. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 07 (2015) 143 DOI: 10.1007/JHEP07 (2015)143 Cite as: arXiv:**

**1505.01106 [hep-lat] (or arXiv: 1505.01106v2 [hep-lat] for this version) Hagedorn spectrum and thermodynamics of SU (2) and SU (3) Yang-Mills theories Michele Caselle, Alessandro Nada, Marco Panero STOP B. de Wit J. Hoppe, H. Nicolai study the quantum-mechanical properties of (e) a supermembrane and examine the nature of (e) its ground state. A supersymmetric gauge theory of (e) area-preserving transformations provides (eb) a convenient framework for this study. The supermembrane can be viewed as (=) a limiting case of a class of models in supersymmetric quantum mechanics. Its mass does not depend on (=) the zero modes and vanishes only if (e) the wave function is (=) a singlet under (e&eb) supersymmetry transformations of (e) the nonzero modes. Authors exhibit (eb) the complexity of (e) the supermembrane ground state and (e&eb) examine various truncations of these models. None of these truncations has massless states. Copyright © 1988 Published by Elsevier B.V. Nuclear Physics B Volume 305, Issue 4, 28 December 1988: Pages 545-581: On the quantum mechanics of supermembranes Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. B. de Wit J. Hoppe, H. Nicolai STOP In this note we explain how world-volume geometries of (e) D-branes can be reconstructed within (eb) the microscopic framework where (e) D-branes are described through (e&eb) boundary conformal field theory. Authors extract the (non-commutative) world-volume algebras from (e) the operator product expansions of (e) open string vertex operators. For branes in a flat background with (e&eb) constant non-vanishing B-field, the operator products are computed perturbatively to (e) all orders in the field strength. Resulting series coincides with (=) Kontsevich's presentation of the Moyal product. After extending these considerations to (e) fermionic fields authors conclude with some remarks on the generalization of approach to (e&eb) curved backgrounds. D-branes and deformation quantization Volker Schomerus1,2 Published 25 July 1999 Journal of High Energy Physics, Volume 1999, JHEP06(1999)STOP Authors employ the path integral approach developed in Gangopadhyay S. and Scholtz F. E., Phys. Rev.**

**Lett., 102 (2009) 241602 to discuss the (generalized) harmonic oscillator in (eb) a noncommutative plane. The action for (e) this system is derived in (e&eb) the coherent state basis with (e&eb) additional degrees of freedom. From this the action in the coherent state basis without (e) any additional degrees of freedom is obtained. This gives (eb) the ground-state spectrum of the system. Authors then employ the exact renormalization group approach to show (eb) that equivalence can be constructed between this (noncommutative) system and (e&eb) a commutative system Path integral action and exact renormalization group dualities for quantum systems in noncommutative plane Sunandan Gangopadhyay Published 25 June 2015 • Copyright © EPLA, 2015 STOP Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini present a lattice study of (e)the equation of state in Yang-Mills theory based on (e) the exceptional G(2) gauge group. As is well-known, at zero temperature this theory shares (e&eb) many qualitative features with real-world QCD, including (e) the absence of colored states in the spectrum and dynamical string breaking at (e) large distances. In agreement with previous works, authors show that at finite temperature this theory features (eb) a first-order deconfining phase transition, whose nature can be studied by (e) a semi-classical computation. Authors also show that the equilibrium thermodynamic observables in the deconfined phase bear striking quantitative similarities with (=) those found in SU (N) gauge theories: in particular, these quantities exhibit (eb) nearly perfect proportionality to the number of gluon degrees of freedom, and (e&eb) the trace anomaly reveals (eb) a characteristic quadratic dependence on (e&eb) the temperature, also observed in SU (N) Yang-Mills theories (both in four and in three spacetime dimensions). Authors compare our lattice data with (e&eb) analytical predictions from (e) effective models, and discuss their implications for (e) the deconfinement mechanism and (e&eb) high-temperature properties of (e) strongly interacting, non-supersymmetric gauge theories. Results give (eb) strong evidence for (e) the conjecture that the thermal deconfining transition is governed by (e&eb) a universal mechanism, common to (e) all simple gauge groups.**

**Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 03 (2015) 057 DOI: 10.1007/JHEP03 (2015)057 Report number: DESY-14-146; IFT-UAM/CSIC-14-076 Cite as: arXiv: 1409.8305 [hep-lat] (or arXiv: 1409.8305v2 [hep-lat] for this version Exceptional thermodynamics: The equation of state of G(2) gauge theory Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini STOP In U (1) lattice gauge theory in three spacetime dimensions, the problem of confinement can be studied analytically in (e&eb) a semi-classical approach, in terms of (e&eb) a gas of monopoles with (e&eb) Coulomb-like interactions. In addition, this theory can be mapped to (e&eb) a spin model via (e&eb) an exact duality transformation, which allows (eb) one to perform high-precision numerical studies of (e) the confining potential. Taking advantage of these properties, authors carried out an accurate investigation of (e) the effective string describing (eb) the low-energy properties of flux tubes in (eb) this confining gauge theory. Authors found striking deviations from (e) the expected Nambu-Goto-like behavior, and, for the first time, evidence for (e) contributions that can be described by a term proportional to the extrinsic curvature of the effective string worldsheet. Such term is allowed by (e) Lorentz invariance and its presence in the infrared regime of the U (1) model was indeed predicted by (e) Polyakov several years ago. Results show that this term scales as expected according to (e) Polyakov's solution, and becomes (=) the dominant contribution to (e) the effective string action in the continuum limit. Authors also demonstrate analytically that the corrections to (e) the confining potential induced by (e) the extrinsic curvature term can be related to (e&eb) the partition function of (e) the massive perturbation of (e) a c=1 bosonic conformal field theory. The implications of results for SU (N) Yang-Mills theories in (e&eb) three and in four spacetime dimensions are discussed. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th) Journal reference: JHEP 1501 (2015) 105 DOI: 10.1007/JHEP01 (2015)105 Report number: IFT-UAM/CSIC-14-052 Cite as: arXiv:**

**1406.5127 [hep-lat] (or arXiv: 1406.5127v2 [hep-lat] for this version) A different kind of string Michele Caselle, Marco Panero, Roberto Pellegrini, Davide Vadacchino STOP In the past ten years, the ideas of supersymmetry have been profitably applied to (e&eb) many nonrelativistic quantum mechanical problems. In particular, there is now a much deeper understanding of why certain potentials are analytically solvable and (e&eb) an array of powerful new approximation methods for (e) handling potentials which are not (e) exactly solvable. In this report, authors review the theoretical formulation of (e) supersymmetric quantum mechanics and discuss many applications. Exactly solvable potentials can be understood in terms of (e&eb) a few basic ideas which include (e) supersymmetric partner potentials (e&eb) shape invariance and (e&eb) operator transformations. Familiar solvable potentials all have (e) the property of shape invariance. Authors describe new exactly solvable shape invariant potentials which include (e) the recently discovered self-similar potentials as (=) a special case. The connection between inverse scattering,(e&eb) isospectral potentials and (e&eb) supersymmetric quantum mechanics is discussed and multisoliton solutions of (e) the K dV equation are constructed. Approximation methods are also discussed within (eb) the framework of supersymmetric quantum mechanics and in particular it is shown that a supersymmetry inspired WKB approximation is (=) exact for a class of shape invariant potentials. Supersymmetry ideas give (eb) particularly nice results for the tunneling rate in (eb) a double well potential and for improving (eb+) large N expansions. Authors also discuss the problem of (e) a charged Dirac particle in (eb) an external magnetic field and other potentials in terms of (e&eb) supersymmetric quantum mechanics. Finally, authors discuss structures more general than supersymmetric quantum mechanics such as(=) parasupersymmetric quantum mechanics in which there is symmetry between a boson and (e&eb) a para-fermion of order p. Copyright © 1995 Published by Elsevier B.V. Physics Reports Volume 251, Issues 5–6, January 1995, Pages 267–385Supersymmetry and quantum mechanics Fred Coopera, Avinash Khareb, Uday Sukhatmec STOP The free energy of (e) the maximally**

**supersymmetric SU (N) gauge theory at temperature T is expected to scale, in the large-N limit, as (=) N2T4 times a function of the 't Hooft coupling, f (gYM2N). In the strong coupling limit the free energy has been deduced from (e) the near-extremal 3-brane geometry, and its normalization has turned out to be Full-size image (<1 K) times that found in (eb) the weak coupling limit. In this paper authors calculate the leading correction to (e) this result in inverse powers of (e) the coupling, which originates from (e) the R4 terms in the tree-level effective action of (e) type IIB string theory. The correction to Full-size image (<1 K) is positive and of (e) order Full-size image (<1 K). Thus, f (gYM2N) increases as (e&eb) the 't Hooft coupling is decreased, in accordance with (e&eb) the expectation that is should be approaching 1 in (eb) the weak coupling limit. Authors also discuss similar corrections for (e) other conformal theories describing (eb) coincident branes. In particular, authors suggest (eb) that the coupling independence of (e) the near extremal entropy for (e) D1-branes bound to (e) D5-branes is related to (e&eb) the vanishing of (e) the Weyl tensor of AdS3 × S3. Copyright © 1998 Published by Elsevier B.V. Nuclear Physics B Volume 534, Issues 1–2, 23 November 1998, Pages 202–222Coupling constant dependence in the thermodynamics of N = 4 supersymmetric Yang-Mills theory Steven S. Gubsera, Igor R. Klebanovb, 1, Arkady A. Tseytlinc, 2 STOP A new approach to (e) N=2 supersymmetry based on (e) the concept of harmonic superspace is proposed and is used to give (eb) an unconstrained superfield geometric description of N=2 super Yang-Mills and supergravity theories as (=) well as of matter N=2 hypermultiplets. The harmonic N=2 superspace has (e) an independent coordinates, in addition to the usual ones, the isospinor harmonics ui+or- on (e&eb) the sphere SU (2)/U (1). The role of ui+or- is to relate the SU (2) group realised on the component fields to (e&eb) a U (1) group acting on (e&eb) the relevant superfields. Their introduction makes it possible to SU (2)-covariantise (eb (t), e&eb) the notion of Grassmann analyticity. Crucial for the construction is the existence of (e) an analytic subspace of the general harmonic N=2 superspace. The hypermultiplet superfields and (e&eb) the true prepotentials (pre-**

**prepotentials) of N=2 super Yang-Mills and supergravity are (=) unconstrained superfunctions over (e&eb) this analytic subspace. The pre-prepotentials have (e) a clear geometric interpretation as gauge connections with (e&eb) respect to the internal SU (2)/U (1) directions. A radically new feature arises: the number of gauge and (e&eb) auxiliary degrees of freedom becomes (=) infinite while (e) the number of physical degrees of freedom remains (=) finite. Other new results are the massive N=2 Yang-Mills theory and (e&eb) various off-shell self-interactions of hypermultiplets. The propagators for matter and Yang-Mills superfields are given. Unconstrained N=2 matter, Yang-Mills and supergravity theories in harmonic superspace A Galperin, E Ivanov, S Kalitzin, V Ogievetsky and E Sokatchev STOP By assuming the existence of a novel multipronged string state for (e) D-particles interacting with (e&eb) D-brane intersections in type IIA string theory, authors are able to derive a quantum mechanical description of (e) supersymmetric Reissner-Nordstrom black holes. A supersymmetric index calculation provides (eb) evidence for (e) this conjecture. The quantum mechanical system becomes (=) two decoupled conformal quantum mechanical systems in (eb) the low energy limit. The conformal quantum mechanics has expected properties of (e) a dual description of string theory on (e&eb) AdS2×S2. Subjects: High Energy Physics - Theory (hep-th) Journal reference: JHEP 0001:006, 2000 DOI: 10.1088/1126-6708/2000/01/006 Report number: LBNL-43734 Cite as: arXiv:hep-th/9907101 (or arXiv:hep-th/9907101v5 for this version) Supersymmetric Quantum Mechanical Description of Four Dimensional Black Holes Julie D. BlumSTOP Using recently proposed soliton equations authors derive a basic identity for the scaling violation of (e) N=2 supersymmetric gauge theories Σiai∂F/∂ai−2F=8πib1u. Here F is (=) the prepotential, ai's are (=) the expectation values of (e) the scalar fields in (eb) the vector multiplet, u=1/2 Tr <ϕ2> and b1 is the coefficient of (e) the one-loop β-function. This equation holds in (eb) the Coulomb branch of all N=2 supersymmetric gauge theories coupled with (e&eb) massless matter. TOHRU EGUCHI and SUNG-KIL YANG, Mod Phys. Lett. A 11, 131 (1996).**

**DOI: 10.1142/S0217732396000151 PREPOTENTIALS OF N=2 SUPERSYMMETRIC GAUGE THEORIES AND SOLITON EQUATIONS STOP The Coulomb branch of N = 2 supersymmetric gauge theories in four dimensions is described in general by (e) an integrable Hamiltonian system in (eb) the holomorphic sense. A natural construction of such systems comes from (e) two-dimensional gauge theory and (e&eb) spectral curves. Starting from this point of view, authors propose an integrable system relevant to (e&eb) the N = 2 SU (n) gauge theory with (e&eb) a hypermultiplet in the adjoint representation, and offer (eb) much evidence that it is correct. The model has (e) a Full-size image (<1 K) S-duality group (with the central element −1 of Full-size image (<1 K) acting as (=) charge conjugation); Full-size image (<1 K) permutes the Higgs, confining, and oblique confining phases in the expected fashion. We also study more exotic phases. Copyright © 1996 Published by Elsevier B.V. Nuclear Physics B Volume 460, Issue 2, 5 February 1996, Pages 299–334Supersymmetric Yang-Mills theory and integrable systems Ron Donagia, 1, Edward Wittenb, 2 STOP A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 present a first step towards generalizing (eb) the work of Seiberg and Witten on N = 2 supersymmetric Yang-Mills theory to (e&eb) arbitrary gauge groups. Specifically, authors propose a particular sequence of (e) hyperelliptic genus n − 1 Riemann surfaces to underly (e&eb) the quantum moduli space of SU (n)N = 2 supersymmetric gauge theory. These curves have (e) an obvious generalization to (e) arbitrary simply laced gauge groups, which involves (e&eb) the A-D-E type simple singularities. To support proposal, authors argue that the monodromy in (e&eb) the semiclassical regime is (=) correctly reproduced. We also give some remarks on a possible relation to string theory. Copyright © 1995 Published by Elsevier B.V. Physics Letters B Volume 344, Issues 1–4, 26 January 1995, Pages 169–175Simple singularities and N = 2 supersymmetric Yang-Mills theory A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 STOP Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 review and extend earlier work that uses (e) the AdS/CFT correspondence to relate the**

**black-hole–black-string transition of gravitational theories on a circle to (e&eb) a phase transition in maximally supersymmetric (1 + 1)-dimensional SU(N) gauge theories at (eb) large N, again compactified on (e&eb) a circle. Authors perform gravity calculations to determine (eb) a likely phase diagram for (e) the strongly coupled gauge theory. Authors then directly study the phase structure of (e) the same gauge theory, now at (eb) weak't Hooft coupling. In the interesting temperature regime for (e) the phase transition, the (1 + 1)-dimensional theory reduces to (e&eb) a (0 + 1)-dimensional bosonic theory, which they solve using (e) Monte Carlo methods. Authors find strong evidence that the weakly coupled gauge theory also exhibits (eb) a black hole–black string-like phase transition in (eb) the large N limit. Authors demonstrate that a simple Landau–Ginzburg-like model describes (eb) the behaviour near the phase transition remarkably well. The weak coupling transition appears to be (=) close to the cusp between a first-order and a second-order transition. Black-hole–black-string phase transitions in thermal (1 + 1)-dimensional supersymmetric Yang–Mills theory on a circle Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 Published 29 October 2004 • 2004 IOP Publishing Ltd Classical and Quantum Gravity, Volume 21, Number 22STOP Monitoring the orbits of stars around (e&eb) Sgr A* offers(eb)the possibility of detecting (e&eb) the precession of (e) their orbital planes due to (e) frame dragging, of (e) measuring the spin and quadrupole moment of (e) the black hole, and of (e)testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and (e&eb) tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are (=), in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for which an optical interferometer, such as GRAVITY, will detect orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to (e&eb) the timescale of precession due to (e) the quadrupole moment of (e) the black hole. As a result, this non-conservative effect is (=) a**

**potential source of systematic uncertainty in (eb) testing the no-hair theorem with (e&eb) stellar orbits. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI:10.1088/0004-637X/777/1/57 Cite as:arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this version) Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP Authors perform a phase space analysis of (e) a generalized modified gravity theory with (e&eb) non-minimal coupling between geometry and (e&eb) matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that (eb) in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, Authors prove that there is (=) an attractor solution to describe (eb) the late time accelerating universe when (e) the modified gravity is chosen in (eb) a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and supernova type Ia data sets, Authors find that (eb) the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as:arXiv: 1512.09281 [gr-qc] (or arXiv: 1512.09281v1 [gr-qc] for this version) Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang study the P−V criticality and phase transition in (eb) the extended phase space of (e) charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of the system and (e&eb) its conjugate quantity is (=) the thermodynamic volume of (e) the black hole. The black**

**holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb)no P−V criticality and phase transition appear, while for the black holes with a spherical horizon, even when (e) the charge of the black hole is absent, the P−V criticality and the small black hole/large black hole phase transition (eb) will appear, but it happens only in (eb) d=5 dimensions; when (e) the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to (e) have the P−V criticality and the small black hole/large black hole phase transition, there exists (eb) an upper bound for (e) the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the Gauss-Bonnet coefficient and Q is the charge of the black hole. Authors calculate the critical exponents at the critical point and find that for all cases, they are (=) the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09 (2013)005 Cite as: arXiv: 1306.6233 [gr-qc] (or arXiv: 1306.6233v4 [gr-qc] for this version) P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP paper studies the whole process of (e) gravitational collapse and (e&eb) accretion of (e) a massless scalar field in (eb) asymptotically flat spacetime. Two kinds of initial configurations are considered. One is the initial data without (e) black hole; the other contains (e) a black hole. Under suitable initial conditions, Authors find that multi-horizon will appear which mean (eb) that the initial black hole formed by (e) gravitational collapse or existing at the beginning will instantly expand and suddenly grow (e&eb, eb), rather than grows gradually in the accretion process. A new type of critical behavior is found around (e&eb) the instant expansion. The numerical computation shows (eb) that the critical exponents are (=) universal. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Computational Physics (physics. Comp-ph) Cite as: arXiv: 1604.03363 [gr-qc] (or arXiv: 1604.03363v1 [gr-qc] for this version)**

**Multi-horizon and Critical Behavior in Gravitational Collapse of Massless Scalar Zhoujian Cao, Rong-Gen CAI, Run-Qiu Yang STOP In this paper, Authors present the results on the gravitational collapse of (e) charged massless scalar field in asymptotically flat spacetime with(e&eb) a perfectly reflecting wall. Differing from previous works, Authors study the system in (eb) the double null coordinates, by which we could simulate (e&eb) the system until the black hole forms with (e&eb) higher precision but less performance time. Authors investigate the influence of charge on (e&eb) the black hole formation and the scaling behavior near the critical collapses. The gapless and gapped critical behaviors for (e) black hole mass and charge are studied numerically. Authors find that they satisfy (eb) the scaling laws for critical gravitational collapse but the gapped critical behavior is different from (e) its AdS counterpart. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1602.00112 [gr-qc] (or arXiv: 1602.00112v1 [gr-qc] for this version) Multiple critical gravitational collapse of charged scalar with reflecting wall Rong-Gen CAI, Run-Qiu Yang STOP In this paper Authors investigate the (2+1)-dimensional topological non-Fermi liquid in (eb) strongly correlated electron system, which has (e) a holographic dual description by (e) Einstein gravity in (3+1)-dimensional anti-de Sitter (AdS) space-time. In a dyonic Reissner-Nordstrom black hole background, Authors consider a Dirac fermion coupled to (e&eb) the background U (1) gauge theory and an (e&eb) intrinsic chiral gauge field bM induced by (e) chiral anomaly. UV retarded Green's function of (e) the charged fermion in (eb) the UV boundary from (e) AdS4 gravity is calculated, by imposing (e&eb) in-falling wave condition at the horizon. Authors also obtain IR correlation function of (e) the charged fermion at (eb) the IR boundary arising from (e) the near horizon geometry of (e) the topological black hole with (e&eb) index k=0,±1. By using (e) the UV retarded Green's function and IR correlation function Authors analyze the low frequency behavior of (e) the topological non-Fermi liquid at zero and (e&eb) finite temperatures, especially the relevant non-Fermi liquid behavior near (eb) the quantum critical region. In addition, Authors**

**find that when k=±1, the effective mass of (e) bulk Dirac fermion at IR fixed point is (=) topologically quantized, consequently the Fermi momentum presents (eb) a discrete sparse distributed pattern, in contrary to (e&eb) the Ricci flat case with k=0. As demonstrated examples, Authors calculate the spectral functions of (e) topological non-Fermi liquid due to (e) anomalous Hall effect (AHE) (|b→|≠0,b 0=0) and chiral magnetic effect (CME) (|b→|=0,b0≠0), which can be tested in topological materials such as (=) topological insulator and (e&eb) topological metal with (e&eb) strong couplings. Subjects: High Energy Physics - Theory (hep-th); Strongly Correlated Electrons (cond-mat.str-el); General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv: 1601.03865 [hep-th] (or arXiv: 1601.03865v1 [hep-th] for this version) Topological Non-Fermi Liquid Rong-Gen Cai, Yong-Hui Qi, Yue-Liang Wu, Yun-Long Zhang STOP The unusual magnetic materials are significant in both science and technology. However, because of (e) the strongly correlated effects, it is difficult to understand their novel properties from (e) theoretical aspects. Holographic duality offers (eb) a new approach to understanding such systems from (e) gravity side. This paper will give a brief review of Authors recent works on the applications of holographic duality in understanding (e&eb) unusual magnetic materials. Some quantitative compare between holographic results and (e&eb) experimental data will be shown and some predictions from (e) holographic duality models will be discussed. Comments: For Second LeCosPA Symposium at Taiwan National University, Taipai, Dec. 14-18, 2015 Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1601.02936 [hep-th] (or arXiv:1601.02936v1 [hep-th] for this version) Understanding strongly coupling magnetism from holographic duality Rong-Gen Cai, Run-Qiu Yang STOP It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a**

**Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3**

**[astro-ph.CO] for this version) Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang**

# Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

## Module Number Six

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-465) 1990. We demonstrate its applicability cosmology/ quantum computation economics etal. For the first time, physicists couple distant nuclear spins using a single electron July 11, 2016. Some sentences have been deleted due to spatial restrictions. Please excuse me. For the first time, researchers at the University of Basel in Switzerland have coupled the nuclear spins of distant atoms using (e) just a single electron. Three research groups took part in this complex experiment, the results of which have now been published in the journal Nature Nanotechnology. In most materials, the nuclear spins of (e) neighboring atoms have (e) only a very weak effect on one another, as (=) the tiny nuclei are located deep within (e&eb) the atoms. This is not the case in metals, however, where (e) some of the electrons are able to move freely. The electron spins are able to couple nuclear spins at (eb) relatively large distances from one another. Named after four physicists, this RKKY interaction was discovered in the 1950s. Nuclear spins linked by (e&eb) an individual electron spin Now, an experiment by researchers from the Department of Physics at the University of Basel has for the first time succeeded in demonstrating this mechanism with (e&eb) a single electron, describing it using (e) quantum theory. The team led by Prof. Richard Warburton trapped a single electron inside (eb) a quantum dot. With use of a method developed in Basel to measure (eb) nuclear spin resonance, they showed that the electron coupled (e&eb) the spins of nuclei up to five nanometers apart - a huge distance in the world of (e) nuclear spins. The results are particularly relevant to (e&eb) the development of spin qubits: these seek to use (e) electron spins to carry (e&eb) information, but the interaction with (e&eb) the nuclei limits (e) the stability of the quantum information. Concentrated blast of physics More information: Role of the electron spin in determining the coherence of the nuclear spins in a quantum dot, Nature Nanotechnology, DOI: 10.1038/nnano.2016.114 Journal reference: Nature Nanotechnology search and more info website Provided by: University of Basel Read more at: http://phys.org/news/2016-07-physicists-couple-distant-nuclear-electron.html#jCp A new spin on quantum computing: Scientists train electrons with microwaves February 15, 2016 A silicon sample doped with (e&eb) bismuth atoms (left image) that is (=) just 150 nanometers thick is fitted with (e&eb) a superconducting resonator that includes (e) a capacitor (black, in left image; light gray in center image) and an inductive wire (red line … In what may provide (eb) a potential path to processing information in (eb) a quantum computer, researchers have switched an intrinsic property of electrons from an excited state to (e&eb) a relaxed state on demand using (e) a device that served as (=) a microwave "tuning fork." The team's findings could also lead to (eb) enhancements in magnetic resonance techniques, which are (=) widely used to explore (e) the structure of materials and biomolecules, and for medical imaging. The**

**international research team, which included scientists at the Department of Energy's Lawrence Berkeley National Laboratory (Berkeley Lab), demonstrated how to dramatically increase the coupling of (e) microwaves in a specially designed (e&eb) superconducting cavity to (e) a fundamental electron property called spin—which, like a coin, can be flipped. By zapping an exotic silicon material developed at Berkeley Lab with the microwaves, they found that they could rapidly change (e&eb) the electron spins from (e) an excited state to a relaxed, ground state by (e) causing (eb) the electrons to emit (eb) some of their energy in the form of (e) microwave particles known as photons. Left on their own, the electron spins would be extremely unlikely to flip back to (e) a relaxed state and to also emit (eb) a photon - the natural rate for this light-emitting effect, known as the Purcell effect, is about once every 10,000 years. The experiment demonstrated an accelerated, controllable relaxation of (e) electron spins and the release of (e) a microwave photon in about 1 second, said Thomas Schenkel, a physicist in Berkeley Lab's Accelerator Technology and Applied Physics Division who led the design and development of the silicon-bismuth sample used in the experiment. A new spin on quantum computing: Scientists train electrons with microwaves Like the classical tuning fork slung from guitar strings in this photo, scientists used (e) a device that worked like a microwave tuning fork to switch (e&eb) a fundamental property of electrons on (e&eb) demand in a silicon-bismuth sample developed at Berkeley Lab. Credit: Flickr/Pierre Guinoiseau "It's like a juggler who throws the balls up, and the balls come down 1,000 times faster than normal, and they also emit (eb) a microwave flash as they drop," he said. The results were published online Feb. 15 in the journal Nature. "Our results are highly significant for (e&eb) quantum information processing," said Patrice Bertet, a quantum electronics scientist at the French Atomic Energy Commission (CEA) who led the experiment. "Indeed, they are a first step toward (e&eb) the strong coupling of (e) individual electron spins to microwave photons, which could form (eb) the basis of new spin-based quantum computer architecture." John Morton, a professor at the London Center for Nanotechnology and co-author of the study, said, "Our ultimate aim is to find a link between quantum information that is fixed and (e&eb) quantum information that can be transported by photons." In today's computers, information is stored as (=) individual bits, and each bit can either be (=) a one or a zero. Quantum computers, though, could conceivably be (=) exponentially more powerful than modern computers because (e) they would use a different kind of bit, called a qubit, which because of (e) the weird ways of quantum mechanics can simultaneously behave (e&eb) as both a one and a zero. A coupled array of qubits would allow (eb) a quantum computer to perform (e&eb) many, many calculations at (eb) the same time, and electron spins are (=) candidates for qubits in (eb) a quantum computer. The latest study shows how the microwave photons could work in concert with (e&eb) the spins of electrons to move information in a new type of computer. "What we need now is ways to wire up these systems—to couple (e&eb) these spins together," Morton said. "We need to make coupled qubits that can perform (e&eb) computations." In the experiment, conducted at CEA in France, a small sample of a highly purified form of (e) silicon was implanted with (e&eb) a matrix of bismuth atoms, and a superconducting aluminum circuit was deposited on (e&eb) top to create (eb) a high-quality resonant cavity that allowed (eb) precise tuning of the microwaves. The electron spins of the bismuth atoms were then flipped into (e&eb) the excited, "spin-up" state. The microwave cavity was then tuned (e&eb) like a musical tuning fork, to (e) a particular resonance that coaxed the spins into emitting (eb) a photon as they flipped back to (e) a relaxed state. The cavity boosted the number of states into (e&eb) which a photon can be (=) emitted, which greatly increased (eb) the decay rate for (e) the electron spins in (eb) a controllable way. The technique is much like buying more lottery tickets to increase (eb) your chances of winning, Morton said. A new spin on quantum computing: Scientists train electrons with microwaves From left are, Berkeley Lab scientists Thomas Schenkel, Qing Ji and Peter Seidl at the NDCX-II (Neutralized Drift Compression Experiment II), which produces (eb) powerful ion beams. Researchers are exploring how to use NDCX-II to process exotic … The large bismuth atoms embedded in (eb) the silicon sample provided (eb) the electrons with (e&eb) unique spin properties that enabled (eb) the experiment. Schenkel said that implanting the bismuth atoms into (e&eb) the delicate silicon framework, a**

**process known as "doping," was "like squeezing bowling balls into (e&eb) a lattice of ping-pong balls." "We did a new trick with silicon. People wouldn't expect you could squeeze (e) anything new out of (e) silicon," Schenkel said. "Now we're looking into further improving (eb) bismuth-doped silicon and into (e&eb) tailoring the spin properties of (e) other materials, and using (e) this experimental technique for these materials." To enhance the performance of materials used (e) in future experiments, Schenkel said it will be necessary to improve (eb) the doping process so it is less damaging to (e) the silicon lattice. Also, the implantation process could be designed to produce (eb) regularly spaced arrays of (e) individual electron spins that would be more useful for (e) quantum computing than a concentrated ensemble of (e) electron spins. "We are now doing experiments on processing (e&eb) this and other materials at higher temperature and pressure with nanosecond ion pulses at NDCX-II, one of the accelerators here at Berkeley Lab," Schenkel said. "There are indications that it will improve (eb) the overall spin quality." Researchers said the latest research could potentially prove useful in boosting (eb) the sensitivity of scientific techniques like nuclear magnetic resonance spectroscopy and dynamic nuclear polarization, useful for a range of experiments, and could also shorten (e) experimental times by manipulating (e&eb) spin properties. "You need a way to reset (e&eb) spins—the ability to cause (eb) them to relax on demand to improve (eb) the rate at which you can repeat an experiment," Morton said. Bertet said it may be possible to further accelerate (eb) the electron-flipping behavior to (e) below 1 millisecond, compared to (e) the 1-second rate in the latest results. "This will then open the way to many new applications," he said. Explore further: Bismuth provides perfect dance partners for quantum computing qubits More information: A. Bienfait et al. Controlling spin relaxation with a cavity, Nature (2016). DOI: 10.1038/nature16944 Nature | Letter Controlling spin relaxation with a cavity A. Bienfait, J. J. Pla, Y. Kubo,X. Zhou, M. Stern, C. C. Lo, C. D. Weis, T. Schenkel, D. Vion,D. Esteve, J. J. L. Morton& P. Berte 531, 74–77 (03 March 2016) doi:10.1038/nature16944 Spontaneous emission of radiation is (=) one of the fundamental mechanisms by which an excited quantum system returns to (e) equilibrium. For spins, however, spontaneous emission is (=) generally negligible compared to (e&eb) other non-radiative relaxation processes because of (e) the weak coupling between the magnetic dipole and (e&eb) the electromagnetic field. In 1946, Purcell realized that the rate of spontaneous emission can be greatly enhanced by (e) placing the quantum system in a resonant cavity. This effect has since been used extensively to control (e&eb) the lifetime of atoms and semiconducting heterostructures coupled to (e&eb) microwave2 or optical3, 4 cavities, and is essential for (e) the realization of high-efficiency single-photon sources5. Here authors report the application of this idea to (e&eb) spins in solids. By coupling donor spins in silicon to (e&eb) a superconducting microwave cavity with (e&eb) a high quality factor and a small mode volume, authors reach the regime in which spontaneous emission constitutes (e) the dominant mechanism of spin relaxation. The relaxation rate is increased by (e) three orders of magnitude as the spins are tuned to (e) the cavity resonance, demonstrating that energy relaxation can be controlled (e&eb) on demand. Results provide a general way to initialize (eb) spin systems into their ground state and therefore have applications in (eb) magnetic resonance and (e&eb) quantum information processing6. They also demonstrate that the coupling between the magnetic dipole of a spin and (e&eb) the electromagnetic field can be enhanced up to (e) the point at which quantum fluctuations have (e) a marked effect on (e&eb) the spin dynamics; as such, they represent (eb) an important step towards the coherent magnetic coupling of (e) individual spins to (e) microwave photons Scientists obtain evidence of many-body localization in a closed quantum system July 8, 2016 by Olivia Meyer-Streng Science Cover Vol 352, Issue 6293, 24 June 2016, to highlight the research published in this issue (see original publication). The cover shows an artist's view of the experiment, illustrating the persistence of the density step of (e) an atomic distribution in an optical lattice. During equilibration ordinary many-body systems lose (e) all information about (e&eb) the initial state. Every morning authors experience an example for (e) this behaviour. Milk poured into a cup of coffee mixes (e&eb) perfectly and after some time it is impossible to say how exactly the two fluids were (=) put together. The same behaviour holds for (e) almost all quantum systems. However,**

**recently a new phenomenon called "many-body localization" has been predicted theoretically, which allows (eb) well insulated quantum systems to (e) preserve memory of the initial state forever. Now a team of scientists around Dr. Christian Groß and Professor Immanuel Bloch (Director at MPQ and Chair of Quantum Optics at LMU Munich), in cooperation with David Huse (Princeton University), has obtained evidence of such a behaviour in a two-dimensional quantum system of cold rubidium atoms trapped in (e&eb) an optical lattice. The scientists observed that – beyond a certain degree of disorder imprinted on (eb) the particle ensemble in the beginning – the system would relax into (e&eb) a steady state still containing (e) detailed microscopic information about its past. "We were able to observe the transition from a thermalized state into (e&eb) a many-body localized phase", Christian Groß points out. "It is the first observation of that kind in a regime that is not accessible with (e&eb) state-of-the-art simulations on classical computers." The experiment is not only of fundamental interest; the results might also lead to (eb) new ways for storing quantum information. Motivated by the foundational problem of how interacting particles behave (e&eb) in a disordered system, in the 1950s the American physicist Philip Warren Anderson discovered the famous localization phenomenon for (e) non-interacting particles, now called "Anderson localization". Here, disorder prevents (e) the particles to move and consequently all transport is (=) stopped. But what happens when disorder comes together with (e&eb) interactions? Will interactions lead to (eb) transport and thermalization, or will the localization persist even at (eb) high energies? So far, there is no theoretical model that faithfully predicts (eb) the evolution of a closed quantum system in more than one dimension under (e&eb) these conditions, although, the possibility for (e)localization has been theoretically suggested. In order to investigate these questions experimentally, strict requirements on controllability and perfect isolation of (e) the system have to be fulfilled. In the experiment described here, ultracold rubidium atoms are loaded into (eb) an optical lattice, a microscopic array of light traps formed by interfering (e&eb) laser beams. The disorder is created by projecting (e&eb) a computer-generated random light pattern onto (e&eb) the optical lattice. As a result of this "disorder", the depth of each microscopic trap in the lattice potential varies from site to (e) site. In fact, Prof. Bloch's group has advanced the experimental tools to such perfection that they can steer the position of (e) the atoms in the artificial light crystal and (e&eb) the interaction between (e&eb) them almost at will. With a high resolution microscope that detects the emitted fluorescence light, the position of each atom can be observed with (e&eb) high precision. Additionally, the atomic density distribution of the initial state can be (=) controlled, and for variable time intervals the evolution of the distribution can be (=) measured with high precision. These tools at hand, non-thermalizing behaviour can be probed in (eb) a conception ally very simple way. Any thermalized state of a closed system reflects (e&eb) the symmetry of its container – like, for example, water spilled into (e&eb) a round bowl would immediately cover (e) the whole bottom. In analogy, the scientists generate (eb) a density step in the initially prepared sample by (e) blowing away half of the atomic distribution with laser radiation. Then they watch how the remaining particles migrate into (e&eb) the empty half. For small disorders, the initially prepared density step is smeared out (e) fast and the initially empty and filled halves become (=) indistinguishable. However, when the measurement is repeated for (e) strong disorder, traces of the initial state remain and the system does not (e) relax to a thermal state even for (e) very long times. "We observe a fairly sharp onset of non-thermalizing behaviour above (e) a critical value", says Christian Groß. "This absence of thermalization is remarkable because (e) thermalization persists in a system of interacting particles, even at (eb) the high energies probed in the experiment." The scientists interpret their observation as (=) the onset of many-body localization in (eb) the atomic system. This is of fundamental interest because it means (eb) the breakdown of equilibrium statistical mechanics. On the other hand, the persistence of initial state information could be used as (=) a source for quantum information technologies. "It should also be emphasized that we obtain these results for a system size that is far beyond (e) numerically accessible scales", says Jae-yoon Choi, postdoc at the experiment. More information: J.-y. Choi et al. exploring the many-body localization transition in two dimensions, Science (2016). DOI: 10.1126/science.aaf8834 Journal reference: Science search and more**

**info website Provided by: Max Planck Society search and more info website Bosons refusing to thermalize in 2D Messy, interacting quantum-mechanical systems are difficult to analyze theoretically. In a single spatial dimension, the calculations are (=) still tractable, and experiments have recently confirmed (eb) the prediction that sufficiently strong disorder can disrupt (e&eb) the transport of interacting particles. In two dimensions, however, the theoretical blueprint is missing. Choi et al. used (e) single-site imaging of cold 87Rb atoms in (eb) an optical lattice to show (eb) that similar localization occurs in (eb) two-dimensional (2D) systems. The study highlights the power of quantum simulation to solve problems that are currently inaccessible to classical computing techniques. Science, this issue p. 1547 Abstract A fundamental assumption in statistical physics is that generic closed quantum many-body systems thermalize under (e&eb) their own dynamics. Recently, the emergence of many-body localized systems has questioned (e) this concept and challenged our understanding of the connection between statistical physics and (e&eb) quantum mechanics. Here authors report on the observation of a many-body localization transition between thermal and (e&eb) localized phases for (e) bosons in a two-dimensional disordered optical lattice. With our single-site–resolved measurements, authors track the relaxation dynamics of (e) an initially prepared out-of-equilibrium density pattern and find strong evidence for (e) a diverging length scale when (e) approaching the localization transition. Experiments represent a demonstration and in-depth characterization of many-body localization in (eb) a regime not accessible with (e&eb) state-of-the-art simulations on classical computers. Pelican instrument provides crucial experimental evidence of unusual quantum state July 8, 2016 Artist's illustration. Credit: Australian Nuclear Science and Technology Organisation (ANSTO) Experiments using inelastic neutron scattering at the Australian Centre for Neutron Scattering have found indications (eb) of a possible new quantum spin state in a novel antiferromagnetic material barium ytterbium zinc oxide (Ba3Yb2Zn5O11) which provides (eb) both a challenge and validation of (e) the third law of thermodynamics. The quantum mechanical information gained from (e) the experiment is crucial not only in understanding the quantum state of the specific material but also provides (eb) better understanding of high temperature superconductivity and ultimately offers (eb) guidance in the design of materials for (e) quantum computing. A research team led by Associate Professor Takatsugu Masuda at the Institute for Solid State Physics at the University of Tokyo with collaborators from ANSTO, Osaka University and Taiwan's National Cheng Kung University, have performed a successful experiment using Pelican, time-of-flight cold neutron spectrometer at ANSTO. The results have just been published in Physical Review B as a rapid Communication and an Editors' Suggestion. "The great success of this challenging experiment relied on (e) the excellent performance of the Pelican instrument, wonderful technical support at ANSTO and strong expertise in the relevant field from A/Prof Masuda's group," said Dr Dehong Yu and Dr Richard Mole, co-authors and Pelican instrument scientists. Pelican instrument provides crucial experimental evidence of unusual quantum state INS spectra measured at 1.5 K (a), 6 K (b), 12 K (c), and 40 K (d) using setup I. (e) INS spectrum measured at 1.5 K using the setup II. (f) Q dependence of the integrated intensity obtained from spectrum in (e). Red solid curve is the …….(deleted due to spatial constraints) Barium ytterbium zinc oxide (Ba3Yb2Zn5O11) is a novel antiferromagnetic model system having (e) a "breathing pyrochlore lattice" consisting of (e) arrays of alternating large and (e&eb) small corner-sharing tetrahedra of ytterbium atoms. Earlier research had suggested (eb) the material was (=) an isotropic spin half (S= ½) system. However, four low energy excitations measured on the sample with the Pelican instrument at a temperature of 1.5 K clearly disagree with (e) this designation and the associated theoretical model based on (e) the isotropic spins having simple interactions. "These observations suggested (eb) a much more complex spin state in the system," said Yu. "The experiment is interesting and important because as you start going very low in temperature, quantum mechanics takes over (e) system rather than thermodynamics, and you observe unusual properties," said Mole. "According to low temperature heat capacity measurements, the material does not (e) have zero entropy or a single ground state until you reach 63 milli Kelvin or -272.94 Celsius, before it starts obeying the third law." Pelican instrument provides crucial experimental evidence of (e) unusual quantum state**

**Breathing pyrochlore lattice formed by (e) A3+ ions as determined by Kimura, Nakatsuju and Kimura in previous work The third law of thermodynamics states that for pure crystals as the temperature of a system approaches absolute zero K, (eb) its entropy or disorder approaches zero. "The research clearly indicates (eb) that below this temperature the material may be in (eb) an exotic state, such as (=) a liquid-like spin state, in which the arrangement of spins behaves (e&eb) like a liquid. In this exotic state, novel phenomena may show up and be detectable; for example, charge and spin separation of electrons may be (=) possible," said Yu. The researchers suggest that further investigations, experimentally and theoretically, are required to ultimately determine the true identity of (e) the ground state of the material. "Discovering the existence of regular tetrahedral spin made it well worth coming all the way to Australia to make use of this neutron experimental apparatus," said Associate Professor Masuda. Many Japanese researchers are using ANSTO's neutron scattering instruments through the peer review system, while their nuclear facilities are offline. The outcome of this particular experiment is also part of the collaboration agreement between ANSTO and the University of Tokyo. More information: T. Haku et al. Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet, Physical Review B (2016). DOI: 10.1103/PhysRevB.93.220407 Journal reference: Physical Review B search and more info website: provided by: Australian Nuclear Science and Technology Organisation (ANSTO) search and more info website Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 T. Haku, K. Kimura, Y. Matsumoto, M. Soda, M. Sera, D. Yu, R. A. Mole, T. Takeuchi, S. Nakatsuji, Y. Kono, T. Sakakibara, L.-J. Chang, and T. Masuda Phys. Rev. B 93, 220407(R) – Published 24 June 2016 Authors study low-energy excitations in (eb) the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 by (e) a combination of inelastic neutron scattering (INS) and (e&eb) thermodynamical property measurements. The INS spectra are quantitatively explained by (e) spin-12 single-tetrahedron model having (e) XXZ anisotropy and Dzyaloshinskii-Moriya interactions. This model has a twofold degeneracy of (e) the lowest-energy state per tetrahedron and well reproduces (eb) the magnetization curve at (eb) 0.5 K and heat capacity above 1.5 K. At lower temperatures, however, authors observe a broad maximum in (eb) the heat capacity around 63 mK, demonstrating (eb) that a unique quantum ground state is selected due to (e) extra perturbations with (e&eb) an energy scale smaller than the instrumental resolution of INS. Revised 13 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.220407 © 2016 American Physical Society Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA (Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor. In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface. As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone. Such planets are all the "right" distance from (e) their star, and hopefully, also have water. But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone. For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement. But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway. But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain. It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of**

**reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life. Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface. The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113 Abstract The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question. Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry. This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry. Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments. Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy. The electronic and (e&eb) structural variability among them is extraordinary. The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations. The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface. The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers. Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp**

## 39. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_6}{dt} = -b_6 Q_6 + g_6(g_6)Q_6$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_6(0) = 0, \quad g_6(P_6) > 0, \quad P_6 \ 0, \qquad \lim_{P_6 \to \infty} g_6(P_6) = (K_6)_6 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_6 + (\rho_a)_6 + D_6(a_6, P_6, Q_6)\rho_6 + 0, \qquad 0 < a_6 \ \infty,$$

$$\rho_6(t, 0) = \int_0^\infty n_6(a_6, P_6, Q_6)\rho_6(t_6, a_6)da_6$$

Where $D_6$ is the mortality function and $n_6$ is the fertility function.

We assume following equations held

$$\left[\frac{dQ_6}{dt}\right]_6 = x_6\left[\frac{dQ_2}{dt}\right]_6 - x_2\left[\frac{dQ_6}{dt}\right]_6$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_6 - x_4\left[\frac{dQ_2}{dt}\right]_6$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_6 - x_6\left[\frac{dQ_3}{dt}\right]_6$$

Where $x_6, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_6$ and $g_6(P_6)Q_6$

Now when you add the term $(b_6Q_6)$ with $g_6(P_6)(Q_6)$ we are infect see is that $(b_6Q_6)$ is credited incrementally part of $g_6(P_6)(Q_6)$ or instant nearly, when subtracted from $\frac{dQ_6}{dt}$ , it means that $g_6(P_6)(Q_6)$ is slowly eaten by $\frac{dQ_6}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_6 = \mu_6(a_6) + (f_6)_6(P_6)_6 + (f_2)_6(Q)_6 \quad (2)$$

$$n_6 = \beta_6(a_6)h_6(Q_6), \quad (3)$$

Where $\mu_6$the age- dependent debit rate whit is $\mu_6(a_6)$ 0 $\forall a > 0$; $(f_6)_6$ is the density- dependent death rate with $f_6(0) = 0$, $(f_6^6)_6(P_6)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_6)_6(P_6) = (k_6)_6$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_6(0) = 0, (f_2^6)_6(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_6(Q_6) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 6, h^6(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_6 + (\rho_a)_6 = -\mu_6(a_6) + (f_6)_6P_6) + \big((f_2)_6(Q_6)\big)(\rho_6)_6 \quad (4)$$

$$\rho_6(t_6, 0) = h_6(Q_6)\int_0^\infty \beta_6(a_6)\rho_6(t, a_6)da, \quad (5)$$

$$\frac{dQ_6}{dt} = -b_6Q_6 + g_6(P_6)Q_6 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_6$ and $Q_6$ of (4-6) satisfy

$$0 \ \rho_6(t, a_6) \ \rho_6^*(t, a_6),$$

$$0 \leq Q_6(t) \leq Q_6^*(t),$$

Where $\rho_6^*(t, a_6)$ and $Q_6^*(t)$, are solutions of

$$(\rho_t^*)_6 + (\rho_a^*)_6 = -(\mu_6(a_6) + (f_6)_6(P_6^*) + (f_2)_6 Q_6^*))\,\rho_6^*, \quad (7)$$

$$\rho_6^*(t, 0) = \int_0^\infty \beta_6(a_6)\,\rho_6^*(t, a_6)\,da_6 \quad (8)$$

$$\frac{dQ_6^*}{dt} = -b_6\,Q_6^* + g_6(P_6^*)Q_6^*, \quad (9)$$

With $P_6^*(t) = \int_0^\infty \rho_6^*(t, a_6)da_6$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_6^*(t)}{dt} = P_6^*(t)\big(r_6 - (f_6)(P_6^*) - (f_2)_6(Q^*)\big), \quad (10)$$

$$\frac{dQ_6^*(t)}{dt} = Q_6^*(t)(-b_6 + g_6(P_6^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_6(a)e^{-\int_0^a \mu_6(\tau)dt}\,e^{-ra}\,da_6 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_6^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^{\infty} \beta_6(a_6) e^{-\int_0^a \mu_6(t)dt}\, da_6 > 1, \qquad (H1)$$

Which guarantees $r_6 > 0$.

*Remark.* If $\beta_6(a) = [\beta_0]_6 a_6 e^{-\alpha_6 a_6}, \mu_6(a_6) \equiv \mu_6$ then the solution of (12) is

$$r_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_6]_6(P_6) = 0, [f_2](Q_6) = \lambda_6 Q_6\ h_6(Q_6) = 1, \qquad and\ \ g_6(P_6) = c_6 P_6$$

$$[\rho_6]_6 + [\rho_a]_6 = -(\mu_6 + \lambda_6 Q_6)\rho_6$$

$$\rho_6(t,0) = \int_0^{\infty} [\beta_0]_6 a_6 e^{-\alpha_6 a_6} \rho_6(t, a_6)\, da_6$$

$$\frac{dQ_6}{dt} = -b_6 Q_6 + c_6 P_6 Q_6$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_6}{dt} = \omega_6 P_6 - \lambda_6 P_6 Q_6$$

$$\frac{dQ_6}{dt} = b_6 Q_6 - c_6 P_6 Q_6$$

Where $\omega_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$

## 40. THE EQUILIBRIUM $(\rho_0)_6(a_6), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_6(a_6), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_6 + (\rho_a)_6 = -\big(\mu_6(a_6) + (f_6)_6(P_6)\big)\rho,$$

$$\rho_6(t,0) = \int_0^\infty \beta_6(a_6)\rho_6(t,a_6)da_6 \qquad (13)$$

The solution of (63) has the form of

$$\rho_6(t,a_6) = Ce_6^{r(t-a_6)} - \int_0^a (\tau)d\tau - \int_0^6 (f_6)_6\big(P_6(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_6(t) = C_6 e^{r_6 t - \int_0^t (f_6)_6(P_6(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_6^{-6}(r)$ if $k_6 > r$. This is equivalent to

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{(k_6)_6 a_6} da_6 < 1.$$

In fact, if we define

$$R_6(s_6) = \int_0^\infty \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau)d\tau} e^{-s_6 a_6}\, da_6$$

$R(s)$ is a decreasing function. $R_6(r_6) = 1$ And $R_6(k_6)_6 < 1$ imply $r_6 < (k_6)_6$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_6(a_6), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{k_6 a_6} da_6$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_6 \equiv \int_0^\infty (\rho_0)_{66}(a_6)da_6$ *satisfies*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{-(f_6)_6(P_0)a_6} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_6(a_6), 0)$

Let $u_6(t, a_6) = \rho(t, a_6) - \rho_0(a_6)$ Then

$$N_6(t) \equiv \int_0^\infty u_6(t, a_6)da_6 = \int_0^\infty \rho_6(t, a_6)da_6 - \int_0^\infty (\rho_0)_6(a_6)da = P_6(t) - P_0)_6$$

And

$$(u_6)_6 + (u_a)_6 = (\rho_t)_6 + (\rho_a)_6 - (\rho_0^6)_6(a_6)$$

$$= -\mu_6(a_6) + (f_6)_6(P_6) + (f_2)_6(Q_6))\,(u_6 + (\rho_0)_6) + \mu_6(a_6)(\rho_0)_6 + (f_6)_6(P_0)_6(\rho_0)_6$$

$$= -\mu_6(a_6)u_6 - (f_6)_6(P_6) + (f_2)_6(Q_6))\,u_6 - (\rho_0)_6\,(f_6)_6(P_6) - (f_6)_6(P_0)_6$$

$$+ (f_2)_6(Q_6)), \qquad (15)$$

Where $(\rho_0^6)_6(a_6)$ denotes $d(\rho_0)_6(a_6)/da_6$

Expansion $(f_6)_6(P_6)$ and $(f_2)_6(Q_6)$ at $(P_0)_6$ and 0 yield respectively

$$(f_6)_6(P_6) = (f_6)_6((P_0)_6 + N_6)$$

$$= (f_6)_6(P_0)_6 + (f_6^6)_6(P_0)_6 N_6 + o_6(N_6), \qquad (16)$$

$$(f_2)_6(Q_6)$$

$$= (f_2^6)_6(0)Q + o_6(Q_6) \quad (17)$$

Write $f_6(P_0) = f_{60}, f_6'(P_0) = f_{60}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_6)_6 + (u_a)_6 \approx -(\mu_6(a_6) + (f_{60})_6)u - (f_{60}^6)_6(\rho_0)_6(N_6) - (f_{20}^6)_6(\rho_0)_6(Q_6)$$

$$u_6(t,0) = \rho_6(t,0)(\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a) + (\rho_0)_6(a_6)da_6 - (\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a_6)da_6 + (h_6(Q_6) - 1)(\rho_0)_6(0) \quad (18)$$

Expanding

$$h_6(Q_6) = 1 + (h^6)_6(0)Q_6 + 0(Q_6) \equiv 1 + h_0'Q_6 + 0Q_6$$

And substituting it into (18), we arrive at

$$u_6(t,0) \approx \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + (h_{a_6}')(\rho_0)_6Q_6$$

The linearization is the following :

$$[u_6] + [u_6] = -(\mu_6(a_6) + [t_{60}]_6u_6 - [\alpha_6]_6(a_6)N_6 - [\alpha_2]_6(a_6)Q_6$$

$$u_6(t,0) = \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + [\xi_6]Q_6 \quad (19)$$

$$\frac{dQ_6}{dt} = -b_6Q_6$$

where

$$[\alpha_6]_6(a_6) = [t_6^6 a]_6[p, o]_6(a_6)d2(a_6) = [f_{20}^6]_6(\rho\, o\,)_6(a_6)\xi = [\mathrm{h\ o}^6]_6[\rho \mathrm{o}]_6(\mathrm{o}) \text{ and } \mathrm{b}_6 = \mathrm{b}_6 - g_6(P_6)$$

If $g_6(P_0) > b$, clearly, $(\rho_0)_6(a_6), 0)$ is unstable and the predator population is persistent. Assume $g_6(P_0)_6 < b$.

Then, $Q_6(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_6 + u_a = -(\mu_6(a_6) + (f_{60})_6)u_6 - (\alpha_6)_6(a_6)N_6$$

$$= -(\mu_6(a_6) + f_{60})u_6 - \alpha_6(a_6)\int_0^\infty u_6(t, a_6)da_6$$

$$u_6(t, 0) = \int_0^\infty \beta_6\,(a_6)u_6(t, a_6)da_6$$

Using the method developed by Cushing [2, 3], we suppose

$$u_6(t, a_6) = \tilde{\mathrm{u}}_6(a_6)^{e^{\gamma(t-a_6)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_6}\, da_6$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_6)$ gives,

$$\tilde{\mathrm{u}}_6(a_6) = e^{-\int_0^a \mu_6(\tau_6)d\tau - (f_{60})_6 a}\left(\tilde{\mathrm{u}}_6(0) - w\int_0^a e^{\int_0^\tau \mu_6(s_6)ds + (f_{60})}\tau_6\,(\alpha_6)_6(\tau_6)e^{\gamma_6\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_6(a_6)e^{-\gamma_6 a_6 - (f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6)d\tau}\right)\tilde{\mathrm{u}}_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a(\alpha_6)_6(\tau_6)e^{-\gamma_6\tau_6}e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds}d\tau_6\,da_6$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}da\right)(\tilde{u}_6(0))_6$$

$$+\left(-6-\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds+y\tau}\alpha_6(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(t_6)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_6\,(a_6)Y_6(\gamma_6,a_6)da_6\right)u_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6=0,$$

(24)

$$\left(\int_0^{\infty}Y_6(\gamma_6,a_6)da_6\right)\tilde{u}_6(0)$$

$$+\left(-1-\int_0^{\infty}Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6\right)w=0.$$

Equqtion (20) has a solution of the form of $u_6(t,a_6)=u_6(a_6)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_6(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) Y^{-6}(\gamma_6, a_6) d\tau_6$. We obtain the characteristic equation

$$f_{60}^{l}{}_6 \left( \int_0^\infty Y_6(\gamma_6, a_6) da_6 \right) \left( \int_0^\infty \beta_6(a_6) Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$+ \left( 1 - \int_0^\infty \beta_6(a_6) Y_6(\gamma_6, a_6) da \right) \left( 1 + (f_{60}^{l})_6 \int_0^\infty Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$= 0. \qquad (25)$$

Since

$$(\rho_0)_6(a_6) = (\rho_0)_6(0) e^{-\int_0^a \mu_6(\tau_6) d\tau_6 - (f_{60})_6 a_6}$$

$$I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) e^{\gamma_6 \tau_6 + (f_{60})_6 \tau + \int_0^\tau \mu_6(s_6) ds_6} d\tau$$

$$= \frac{(\rho_0)_6(0)}{\gamma_6} (e^{\gamma_6 a_6} - 1)$$

Let

$$E_6(a_6) = e^{-(f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6) d\tau_6}$$

This relationship is cardinal one. It is useful in quantum computation and other black hole dynamics problems

**Quantum space and quantum time**

(2195) **Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.**

(2196) **In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.**

(2197) **They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.**

(2198) **When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.**

(2199) **Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.**

## The Planck scale

(2200) **Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".**

(2201) **Here ħ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.**

(2202) **The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.**

(2203) **Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)**

(2204) **Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.**

## Physical applications

**Some phenomena where quantum gravity plays a major role are the following.**

### The microscopic structure of spacetime

(2205) **If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.**

(2206) **Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length**

(2207) **Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.**

(2208) **A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.**

**Early cosmology**

(2209) **According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.**

(2210) **However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.**

(2211) **A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.**

(2212) **Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".**

(2213) **One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.**

**Black holes**

(2214) **Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.**

(2215) **First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.**

(2216) **Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.**

(2217) **What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.**

(2218) **Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.**

**Astrophysical effects**

(2219) **Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.**

(2220) **An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.**

(2221) **The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.**

## Main approaches

(2222) **The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.**

(2223) **Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.**

(2224) **The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.**

## String theory

(2225) **String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.**

(2226) **Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.**

(2227) **String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.**

(2228) **Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.**

(2229) **However, when summed the entire series appears to be (=) divergent.**

(2230) **In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.**

(2231) **Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.**

(2232) **Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.**

(2233) **In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.**

**Loop quantum gravity and spin foams**

(2234) **Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.**

(2235) **The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity.**

(2236) **Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime.**

(2237) **The granular properties of space can be (=) explicitly computed.**

(2238) **In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory.**

(2239) **The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks.**

(2240) **Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure.**

(2241) **The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit.**

(2242) **Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism.**

(2243) **This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism.**

**Noncommutative geometry**

(2244) **Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime.**

(2245) **Quantum mechanics teaches us that dynamical quantities are (=) noncommutative -- meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed.**

(2246) **It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry.**

(2247) **A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces.**

(2248) **Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces.**

(2249) **This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity.**

**Others**

(2250) **Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays**

(2251) **. Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection.**

(2252) **Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible.**

(2253) **The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.**

**Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016**

(2254) **These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field.**

(2255) **The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets.**

(2256) **Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state.**

(2257) **Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands.**

(2258) **The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI**

(2259) **The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9].**

(2260) **This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive.**

(2261) **More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.**

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

(2262) **Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.**

(2263) **Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.**

(2264) **The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.**

(2265) **This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.**

(2266) **The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.**

(2267) **This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.**

(2268) **For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.**

(2269) **The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society**

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(2270) **Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.**

(2271) **We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.**

(2272) **When this is true—as is often the case—the system serves as (=) a "bath" for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called thermalization.**

(2273) **However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.**

(2274) **Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1).**

(2275) **Bringing them to thermal equilibrium requires (e) an external bath.**

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| (2276) **But so far, there has been no rigorous proof of (e) the existence of such many-body localization. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption.** |
| (2277) **The concept of localization was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations).** |
| (2278) **The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century.** |
| (2279) **But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information.** |
| (2280) **Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment.** |
| (2281) **For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood.** |
| (2282) **It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7].** |
| (2283) **Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature.** |
| (2284) **One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.** |
| (2285) **However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].** |
| (2286) **Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.** |
| (2287) **Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.** |
| (2288) **As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked.** |
| (2289) **His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice.** |
| (2290) **An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site.** |
| (2291) **This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself).** |

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| (2292) **This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition.** |
| (2293) **However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime:** |
| (2294) **These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects.** |
| (2295) **It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies).** |
| (2296) **This is a realistic scenario for (e) all known systems.** |
| (2297) **On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects.** |
| (2298) **On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other.** |
| (2299) **No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true.** |
| (2300) **His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found.** |
| (2301) **Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization?** |
| (2302) **Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects?** |
| (2303) **Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions.** |
| (2304) **So in those systems, the quantum phase transition does indeed exist.** |
| (2305) **But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension.** |
| (2306) **Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters.** |

**Focus: Biological Cells Form Electric Circuits July 1, 2016• Physics 9, 75 Cells that are electrically active and that also produce light for easy voltage monitoring could lead to new studies of heart arrhythmias and possibly bio-computing. Figure caption expand figure H. M. McNamara et al., Phys. Rev. X (2016) Cell circuit. Waves of electrical activity propagate in a loop within a collection of cells engineered with electrical and optical properties (see video below).Some sentences, pictures and videos are deleted due to constraints on space.**

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| (2307) **The human heartbeat is produced by (e) electrical pulses that propagate through (e&eb) cardiac tissue, causing (eb) rhythmic muscle contraction.** |
| (2308) **Researchers have previously engineered cells to create (eb) an artificial tissue capable of producing (eb) coordinated electrical activity, and now a team has added the ability to monitor (e&eb) their electrical state by (e) detecting fluorescent emission.** |
| (2309) **They have also fashioned the cells into (e&eb) "living circuits" that might act as (=) model systems for studying heart behavior.** |
| (2310) **Constructing electrically active tissues from their basic building blocks by (e) genetically engineering cells allows (eb) researchers to study a "clean" system similar to (e) the more complicated cardiac tissue.** |
| (2311) **These systems let them investigate, for example, the electrical signals that lead to (eb) the potentially fatal irregular heartbeat called arrhythmia.** |
| (2312) **Previously researchers have introduced (eb) ion channels—proteins that conduct ions through (e&eb) cell membranes—into human kidney cells to make (eb) electrically "excitable" tissues [1].** |
| (2313) **Adam E. Cohen of Harvard University and his co-workers decided to add more ingredients to such cells, so that the electrical activity could be (=) light-activated and easily visualized.** |
| (2314) **They incorporated four additional gene products into (e&eb) human embryonic kidney (HEK) cells.** |
| (2315) **Two were ordinary ion channels; the other two were (=) a blue-light-activated ion channel (a light-sensitive switch) and (e&eb) a protein that fluoresces with (e&eb) an intensity that depends on (=) the voltage across the membrane in which it is embedded.** |
| (2316) **Light activation of electrically excitable cells has been shown before [2]. But Cohen and colleagues have now included a convenient "readout" mechanism (the fluorescent protein) and have used patterning methods to shape (e&eb) the cell cultures into "tissue circuits."** |
| (2317) **They grew rings of cells a few millimeters in diameter that could act as (e&eb) oscillators that support (eb) circulating waves of electrical excitation.** |
| (2318) **These waves arise from (e) the collective behavior of the cells: a pulse produced by (e) one cell can trigger another from (e) its neighbors, eventually producing (eb) a coherent wave of excitation traveling through (e&eb) the tissue.** |
| (2319) **It's this coordination, in a repeating cycle of excitation and quiescence, which gives rise to (eb) rhythmic activity in cardiac tissue.** |

**H. M. McNamara et al., Phys. Rev. X (2016)**

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| (2320) **A pattern of blue light (not visible) illuminates these cells, defining (eb) a rectangular loop around which electrical waves can propagate.** |
| (2321) **At about 0:07 the rectangle begins to shrink in (eb) the vertical direction, causing t(eb) he cycle time to gradually decrease. (In other experiments, the cells were grown in (eb) fixed, ring-shaped patterns.)** |
| (2322) **The ring of engineered HEK cells acted like an electronic "ring oscillator," a device with a characteristic frequency to (e&eb) its circulating electrical signal.** |
| (2323) **The electrical waves in a ring oscillator can propagate (e&eb) in either direction around the ring, enabling (eb) it to serve as a binary memory unit.** |

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| (2324) **Cohen and colleagues found that in their HEK tissues, these circulating waves remained (eb) stable for more than two hours, with (e&eb) the pulses completing a few circuits every second.** |
| (2325) **The electrical signals were easily visualized as (=) waves of fluorescent emission.** |
| (2326) **The researchers say that these tissues could be used to (e) understand and predict the behavior of real cardiac cells, especially arrhythmias.** |
| (2327) **"In real heart, if you pace the tissue fast enough, *(eb) it starts skipping beats and eventually becomes (=) arrhythmic," says Cohen.** |
| (2328) **This state is a well-known precursor to (e) sudden cardiac death and is associated with (e&eb) the breakup of the usual regular waves of electrical excitation.** |
| (2329) **Some drugs also have the dangerous side effect of inducing (eb) arrhythmia, and Cohen says that the new system might offer a quick and convenient way of screening for such hazards. The work "is really very elegant and deserves attention," says physiologist Gil Bub of Oxford University in the UK. "It opens new research questions for people like me who are interested in wave dynamics in excitable systems."** |
| (2330) **There might also be practical uses for information processing in living tissues. "By combining patterns of cell growth and (e&eb) patterns of ion channel expression, we think we can build diodes, amplifiers, and other simple logic elements," says Cohen.** |
| (2331) **Andrew Adamatzky of the University of the West of England in Bristol, a specialist in "unconventional computing," foresees circuits that perform computations through (e&eb) the interactions of electrical waves. "This research might lead to fabrication of implantable bio-computers," he says. This research is published in Physical Review X. –Philip Ball Philip Ball is a freelance science writer in London and author of Invisible: The Dangerous Allure of the Unseen (University of Chicago Press, 2014).** |

**Optically Controlled Oscillators in an Engineered Bioelectric Tissue Harold M. McNamara, Hongkang Zhang, Christopher A. Werley, and Adam E. Cohen Phys. Rev. X 6, 031001 – Published 1 July 2016**

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| (2332) **Complex electrical dynamics in excitable tissues occur throughout (e&eb) biology, but the roles of individual ion channels can be difficult to determine due to (e) the complex nonlinear interactions in (eb) native tissue.** |
| (2333) **Here, authors ask whether they can engineer a tissue capable of (e) basic information storage and processing, where (e) all functional components are known and well understood.** |
| (2334) **Authors develop a cell line with (e&eb) four transgenic components: two to enable (eb) collective propagation of electrical waves and (e&eb) two to enable optical perturbation and (e&eb) optical readout of membrane potential.** |
| (2335) **They pattern the cell growth to define (eb) simple cellular ring oscillators that run stably for >2 h (∼104 cycles) and that can store data encoded in (eb) the direction of electrical circulation.** |
| (2336) **Using (e) patterned optogenetic stimulation authors probe the biophysical attributes of (e) this synthetic excitable tissue in detail, including (e) dispersion relations, (e&eb) curvature-dependent wave front propagation, (e&eb) electrotonic coupling, and (e&eb) boundary effects.** |
| (2337) **Authors then apply the biophysical characterization to (e) develop an optically reconfigurable bioelectric oscillator.** |
| (2338) **These results demonstrate the feasibility of engineering bioelectric tissues capable of (e&eb) complex information processing with optical input and output. Received 12 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevX.6.031001 This article is available under the terms of the** |

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**Scientists detect radio emission from a nearby brown dwarf July 6, 2016 by Tomasz Nowakowski report**

(2339) **K2 and Spitzer images of W0607+24: The position of W0607+24 is marked by (e) a red circle. SDSS J060738.06+242939.5, marked by (e) an orange circle, are the brightest star in the K2 image.**

(2340) **The three faint (r = 18.3 − 19.6) stars marked by cyan circles, contribute significant flux to (e) the K2 measurement but are (-) negligible at Spitzer's longer wavelength and (e &êb) higher resolution. Other stars with r < 20 in SDSS are marked with blue circles. Credit: Gizis et al., 2016.**

(2341) **(Phys.org)—Located some 23 light years away, the brown dwarf designated WISEP J060738.65+242953.4 or W0607+24 for short, turns out to be (=) a source of radio emission.**

(2342) **According to a research paper published July 4, this substellar object showcases (eb) quiescent radio emission, making it one of (e) the most radio-faint, ultra-cool dwarfs yet detected.**

(2343) **Brown dwarfs like W0607+24 are objects that are too large to be called planets and too small to be stars. With a mass below that necessary to maintain hydrogen-burning nuclear fusion reactions, they are much cooler and dimmer than (eb) main sequence stars.**

(2344) **These objects could be also sources of (e) radio emission, but the cause of this process is still not completely understood.**

(2345) **Due to its proximity, W0607+24 is (=) an excellent target for studying the nature of (e) brown dwarfs.**

(2346) **It is the nearest known late-L dwarf in the northern hemisphere, and the third-nearest in (=) the whole sky, (eb) allowing researchers to study phenomena occurring in brown dwarfs.**

(2347) **That's why a team of astronomers, led by John Gizis of the University of (e)Delaware, Newark, carried out a series of mid-infrared, radio, and optical observations of W0607+24 to study any activity in this curious substellar object.**

(2348) **Motivated by the results from NASA's prolonged Kepler mission, named K2, which obtained (eb) long time-series photometry of W0607+24; they observed this object simultaneously with the Karl G. Jansky Very Large Array (VLA) in New Mexico and NASA's Spitzer Space Telescope.**

(2349) **The most important results were provided by VLA, as the team found that W0607+24 is (=) a source of quiescent, non-bursting, radio emission.**

(2350) **The scientists noted that the discovery of this process in this nearby object could be (=) crucial for our understanding of this process in (eb) ultracool dwarfs in general.**

(2351) **"W0607+24 is a valuable test case for (e) understanding the mechanisms of radio emission in ultracool dwarfs.**

(2352) **Some of the radio-detected objects emit (eb) rapid, intense bursts with high circular polarization, often found to occur periodically and attributed to coherent emission due to (e) the electron cyclotron maser instability," the paper reads.**

(2353) **The astronomers also suggested that the relative radio faintness of W0607+24 may be due to (e) its orientation if (e) the emission is concentrated at the poles.**

(2354) **However, follow-up observations are needed to better characterize (eb) the detected radio emission and to eliminate (e) all the uncertainties.**

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| (2355) **The team also hopes for more significant findings when (e) studying this nearby substellar object in the future.** |
| (2356) **"As the nearest known northern hemisphere late-L dwarf, this offers many possibilities for future studies. For example, it could serve as (=) an useful standard for (e) line-broadening and (e&eb) polarization studies.** |
| (2357) **If there is a planetary system aligned with (e&eb) (the primary's rotation axis, then we would expect very little radial velocity signal but a strong astrometric signal," the researchers concluded.** |
| (2358) **Moreover, the team plans to monitor (e&eb) additional late-M and L dwarfs in other K2 campaigns, expecting even more promising results.** |
| (2359) **They are more distant than W0607+24, but because they are warmer, (eb) some are brighter and higher signal-to-noise than that described in this paper.** |
| (2360) **In addition to detecting the radio emission, (eb)the scientists also estimated the age, mass and radius of W0607+24** |
| (2361) **. According to the study, the object is less than two billion years old, has (e) a maximum mass of about 0.055 solar masses and a radius approximately of 0.1 solar radii.** |
| (2362) **More information: WISEP J060738.65+242953.4: A Nearby. Pole-On L8 Brown Dwarf with Radio Emission, arXiv: 1607.00943 [astro-ph.SR] arxiv.org/abs/1607.00943** |
| (2363) **Authors present a simultaneous, multi-wavelength campaign targeting the nearby (7.2 pc) L8/L9 (optical/near-infrared) dwarf WISEP J060738.65+242953.4 in (eb) the mid-infrared, radio, and optical.** |
| (2364) **Spitzer Space Telescope observations show no (e) variability at the 0.2% level over 10 hours each in the 3.6 and 4.5 micron bands.** |
| (2365) **Kepler K2 monitoring over 36 days in Campaign 0 rules out (e) stable periodic signals in (eb) the optical with amplitudes great than 1.5% and periods between 1.5 hours and 2 days.** |
| (2366) **Non-simultaneous Gemini optical spectroscopy detects lithium, constraining (e) this L dwarf to be less than ~2 Gyr old, but no Balmer emission is observed.** |
| (2367) **The low measured projected rotation velocity (v sin i < 6 km/s) and lack of variability are (=) very unusual compared to other brown dwarfs, and authors argue that this substellar object is (=) likely viewed pole-on.** |
| (2368) **Authors detect quiescent (non-bursting) radio emission with (e&eb) the VLA.** |
| (2369) **Amongst radio detected L and T dwarfs, it has (e) the lowest observed L_nu and the lowest v sin i.** |
| (2370) **They discuss the implications of pole-on detection for (e) various proposed radio emission scenarios. © 2016 Phys.org** |

**A new look at the galaxy-shaping power of black holes July 6, 2016**

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| (2371) **It wasn't until the advent of X-ray astronomy that the full picture of (e) the galaxy-shaping power of black holes began to emerge (eb) with the ability to see plasma.** |
| (2372) **Data from a now-defunct X-ray satellite is providing (eb) new insights into the complex tug-of-war between galaxies, the hot plasma that surrounds them, and the giant black holes that lurk in their centres.** |

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| (2373) | **Launched from Japan on February 17, 2016, the Japanese space agency (JAXA) Hitomi X-ray Observatory functioned for just over a month before contact was lost and (e&eb) the craft disintegrated.** |
| (2374) | **But the data obtained during those few weeks was enough to paint (eb) a startling new picture of the dynamic forces at work within galaxies.** |
| (2375) | **New research, published in the journal Nature today, reveals data that shows just how important the giant black holes in galactic centres are to the evolution (eb) of the galaxies as a whole.** |
| (2376) | **"We think that supermassive black holes act(e&eb) like thermostats," said Brian McNamara, University Research Chair in Astrophysics at the University of Waterloo. "They regulate (e&eb) the growth of galaxies."** |

**Champagne bubbles of plasma**

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| (2377) | **During its brief life, the Hitomi satellite collected X-ray data from (e) the core of the Perseus cluster, an enormous gravitationally-bound grouping of (e) hundreds of galaxies.** |
| (2378) | **Located some 240 million light years from earth, the Perseus cluster is one of (e) the largest known structures in the universe.** |
| (2379) | **The cluster includes not only the ordinary matter that makes up the galaxies, but an (e) "atmosphere" of hot plasma with a temperature of tens of millions of degrees, as well as a halo of (e) invisible dark matter.** |
| (2380) | **Earlier studies, going back to the 1960s, have shown that each of the galaxies in the cluster - and indeed most galaxies - likely contains (e) a supermassive black hole in its centre, an object 100 million to more than ten billion times as massive as our sun** |

**New look at the galaxy-shaping power of black holes**

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| (2381) | **Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving.** |
| (2382) | **"These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science.** |
| (2383) | **"Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy."** |
| (2384) | **This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return.** |
| (2385) | **The remaining matter gets absorbed into the black hole, adding to (e) its mass.** |
| (2386) | **The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne.** |
| (2387) | **The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution.** |
| (2388) | **Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The** |

**answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form.**

(2389) **"Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara.**

(2390) **"It creates (eb) these bubbles, and the bubbles keep the plasma hot.**

(2391) **That's what prevents galaxies from (e) becoming even bigger than they are now."**

(2392) **Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.**

(2393) **In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.**

(2394) **Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.**

(2395) **"It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.**

(2396) **"It's like nature's thermostat, which keeps these galaxies from (e) growing.**

(2397) **If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."**

**A new look at the galaxy-shaping power of black holes**

(2398) **An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the**

(2399) **McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.**

(2400) **Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.**

**Doomed satellite missions**

(2401) **Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.**

(2402) **At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.**

(2403) **Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a**

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| **leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.** |
| (2404) **The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.** |
| (2405) **These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.** |
| (2406) **"But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people** |
| (2407) **. It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."** |
| (2408) **The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.** |
| (2409) **"This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website** |

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

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| (2410) **Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.** |
| (2411) **They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.** |
| (2412) **However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.** |
| (2413) **Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.** |
| (2414) **X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well.** |
| (2415) **The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma.** |
| (2416) **These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6.** |
| (2417) **Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus.** |

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| (2418) | **A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate.** |
| (2419) | **Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure.** |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (2420) | **This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy.** |
| (2421) | **Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole.** |
| (2422) | **But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy.** |
| (2423) | **A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole.** |
| (2424) | **It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth.** |
| (2425) | **"There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA).** |
| (2426) | **This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies.** |
| (2427) | **However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated.** |
| (2428) | **In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1.** |
| (2429) | **We know dark matter exists only from (e) its gravitational effects.** |
| (2430) | **It holds together galaxies and (e&eb) galaxy clusters.** |
| (2431) | **Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years.** |
| (2432) | **To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies.** |
| (2433) | **They used star motions as tracers to (e) weigh the galaxies' central black holes.** |
| (2434) | **X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto.** |
| (2435) | **They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone.** |
| (2436) | **This connection is likely to be related to (e&eb) how elliptical galaxies grow.** |

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| (2437) **An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together.** |
| (2438) **Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole.** |
| (2439) **"In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp** |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (2440) **The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the …** |
| (2441) **Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness?** |
| (2442) **Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness.** |
| (2443) **"Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain."** |
| (2444) **Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness.** |
| (2445) **"With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication."** |
| (2446) **Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain** |

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| **changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories.** |
| (2447) **The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen.** |
| (2448) **In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer.** |
| (2449) **Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants.** |
| (2450) **These were treated as (=) "aware" and "unaware" trials, respectively.** |
| (2451) **Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it.** |
| (2452) **The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another.** |
| (2453) **Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion.** |
| (2454) **No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness.** |
| (2455) **"We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said.** |
| (2456) **"Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness."** |
| (2457) **The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner.** |
| (2458) **Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain.** |
| (2459) **"We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said.** |
| (2460) **"This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients** |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

(2461) **Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter.**

(2462) **As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined.**

(2463) **Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions.**

(2464) **"Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.**

(2465) **The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.**

(2466) **However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.**

(2467) **The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.**

(2468) **Noda and his colleagues instead studied the change in emissions over (e&eb) time.**

(2469) **By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.**

(2470) **Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).**

(2471) **The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.**

(2472) **Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.**

(2473) **"The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp**

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(2474) **The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking "palm tree" organization.**

(2475) **The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society**

(2476) **(Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible?**

(2477) **With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem.**

(2478) **This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other.**

(2479) **According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible.**

(2480) **This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges.**

(2481) **In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century.**

(2482) **Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching.**

(2483) **In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom.**

(2484) **Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth.**

**Spheres on trees**

(2485) **Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration.**

(2486) **They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement.**

(2487) **Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum.**

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| (2488) **By visualizing the problem in this way, the researchers noticed that these particular graphs don't (e) have lots of branches extending from lots of (e) other branches.** |
| (2489) **Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure.** |
| (2490) **The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network.** |
| (2491) **As a result, it provides (eb) an important clue for finding the global minimum.** |
| (2492) **It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum.** |
| (2493) **This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large.** |

**Thompson's problem 2**

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| (2494) **By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society** |
| (2495) **The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima.** |
| (2496) **The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node.** |

**Implications for Thomson's and Smale's problems**

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| (2497) **Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before.** |
| (2498) **"Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around."** |
| (2499) **As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem.** |
| (2500) **"I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view."** |
| (2501) **In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges.** |
| (2502) **From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear.** |

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| (2503) | **"While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK.** |
| (2504) | **"For larger numbers, there will likely be (=) multiple funnels.** |
| (2505) | **Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum."** |
| (2506) | **Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.** |
| (2507) | **"The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.** |
| (2508) | **"Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org** |

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**Introduction:**

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price, University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss.**

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| **(385)** | **Black Holes are regions of (e) space in which gravitational fields are so strong that no particle or signal can escape (eb) the pull of gravity.** |
| **(386)** | **The boundary of this no-escape region is (=) called the event horizon, since distant observers outside the black hole cannot (e) see (cannot get light from) events inside.** |
| **(387)** | **Although the fundamental possibility of such an object exists within (eb) Newton's classical theory of gravitation, Einstein's theory of gravity makes (eb)black holes inevitable under some circumstances.** |
| **(388)** | **Prior to the early 1960s, black holes seemed to be (=) only an interesting theoretical concept with no (e) astrophysical plausibility, but with (e&eb) the discovery of quasars in 1963 it became clear that very exotic astrophysical objects could exist.** |
| **(389)** | **Nowadays it is taken for granted that black holes do exist in (eb) at least two different forms.** |
| **(390)** | **Stellar mass black holes are (=) the endpoint of the death of some stars, and supermassive black holes are the result of (e) coalescences in the centers of most galaxies, including our own.** |
| **(391)** | **No signal can propagate from (e) inside a black hole, but the gravitational influence of a black hole is (=) always present.** |

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| (392) | This influence does not (e) propagate out of the hole; it is permanently present outside, and depends only on (=) the total amount of mass, angular momentum, and electric charge that have gone into forming (eb) the hole. |
| (393) | Black holes can be detected through the influence of this (e&eb) strong gravity on the surroundings just outside the hole. |
| (394) | In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them. |
| (395) | Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors. |

**Classical vs. relativistic black holes**

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| (396) | Something like a black hole exists within (eb) Newton's classical theory of gravity. |
| (397) | In that theory, an energy argument tells us that there is (=) an escape velocity $v_{esc}=\sqrt{2GM/R}$ from (e) the surface of any spherical object of (e) mass M and radius R |
| (398) | If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity. |
| (399) | Thus the condition for such an "unseeable" object is $R<2GM/c^2$ |
| (400) | In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape. |
| (401) | This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity. |
| (402) | To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it. |
| (403) | The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c . |
| (404) | This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed. |
| (405) | An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity. |
| (406) | Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable. |

**Stationary black holes**

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| (407) | In the Newtonian theory, gravity is described by (e) the potential $\Phi$. |
| (408) | Inside a spherical object the form of (e) $\Phi(r)$ depends on (=) the interior structure, but in the vacuum outside (e) matter the potential $-GM/r$ depends only on (=) the interior mass. |
| (409) | Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object. |

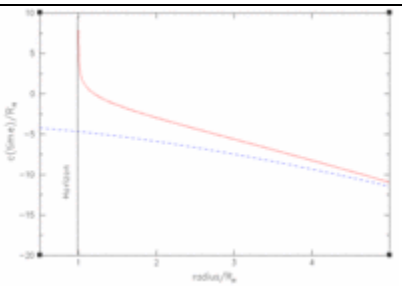
 

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| **(410) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward.** |
| **(411) This horizon radius RH=2GM/c2 is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory.** |
| **(412) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution.** |
| **(413) Radius RHhere actually means that the area of the event horizon is (=) 4πR2H.** |
| **(414) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution.** |
| **(415) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity.** |
| **(416) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own.** |
| **(417) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole.** |
| **(418) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density.** |
| **(419) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed.** |
| **(420) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole.** |
| **(421) In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter.** |
| **(422) Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum.** |
| **(423) In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior.** |
| **(424) The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object.** |
| **(425) Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum.** |
| **(426) These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.]** |
| **(427) All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time.** |
| **(428) But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?"** |
| **(429) The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point.** |

**(430)** One of these coordinates is (=) called the "coordinate time."

**(431)** Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate.



**(432)** Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole

**(433)** At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole.

**(434)** Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations.

**(435)** For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree

**(436)** An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events.

**(437)** The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving.

**(438)** Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle.

**(439)** The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return.

**(440)** But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon.

**(441)** Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future.

**(442)** The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?"

**Black hole parameters**

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| **(443) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin).** |
| **(444) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun.** |
| **(445) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits.** |
| **(446) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form.** |
| **(447) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds.** |
| **(448) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2.** |
| **(449) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed.** |
| **(450) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer.** |
| **(451) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2)** |
| **(452) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials.** |
| **(453) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable** |
| **(454) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite.** |
| **(455) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole.** |
| **(456) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars.** |
| **(457) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity.** |
| **(458) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity.** |
| **(459) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite.** |
| **(460) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity.** |
| **(461) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum.** |

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| **(462) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship.** |
| **(463) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= (4πG2/c4)[(M+M2−a2−−−−−−−√)2+a2].** |
| **(464) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored.** |

**Dynamical black holes**

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| **(465) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape.** |
| **(466) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, a horizon can be born.** |
| **(467) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid.** |
| **(468) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward.** |
| **(469) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations.** |
| **(470) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes.** |
| **(471) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3.** |
| **(472) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/** |
| **(473) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers.** |
| **(474) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission.** |
| **(475) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.** |
| **(476) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.** |
| **(477) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.** |
| **(478) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.** |

**(479)** **The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature 6×10−8 (Mo/M) K.**

**(480)** **Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.**

Then $[Y_6]_6[\gamma_6]_6, a_6) = E_6(a_6)e^{-\gamma_6 a_6}$. by a simple calculation,

$$\int_0^\infty \beta_6(a_6)Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma}\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}\ da\right);$$

$$\int_0^\infty Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma_6}\left(\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}\ da_6\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_6(0)}{\gamma_6}[f_{60}^6]\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}da_6\ (\int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6)$$

$$+\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right).$$

$$\left(1 + \frac{(\rho_0)_6(0)}{\gamma_6}(f_{60}^{'})_6\left(\int_0^\infty E_6(a_6)da_6 - \int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}da_6\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right)(1 + \frac{(f_{60}^{'})_6}{\gamma_6}(\rho_0)_6(0)\int_0^\infty E_6(a_6)da_6) = 0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1 - \int_0^\infty \beta_6(a_6) Y_6(\gamma_6 a_6) da_6 = 1 - \int_0^\infty \beta_6(a_6) E_6(a_6) e^{-\gamma_6 a_6} da_6 \neq 0,$$

Since $\int_0^\infty \beta_6(a_6) Y_6(\gamma_6 a_6) \int_0^a \alpha_6(\tau_6) Y^{-6}(\gamma_6, \tau_6) d\tau_6 da_6 > 0$. Therefore, the characteristic solution is

$$\gamma_6 = -(f_{60}^{'})_6 (\rho_0)_6 \int_0^\infty E_6(a_6) da_6.$$

Since $f_6'(P_0)_6 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_6), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE OF DEBIT TYPES AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_6), 0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_6 + (\mu_6(a_6) + [f_6]_6(P_6^*) + [f_2]_6(Q_6^*))P_6^* = 0, \quad (26)$$

$$p_6^*(0) = h_6(Q_6^*) \int_0^\infty \beta_6(a_6) p_6^*(a_J) p_6^*(a_6) da_6 \quad (27)$$

$g_6(P_6^*) - b_6 = 0,$ (28)

Where $P_6^*=\int_0^\infty p_6^*(a_6)\, da_6$

Equation (28) yields $P^* = g^{-6}(b)$,

From (26),

$p_6^*(a) = p_6^*(0)e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6}$

Substituting into (27), we need to solve

$$\left(1 - h_6(Q_6^*)\int_0^\infty \beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6} da_6\right) p_6^*(0) = 0$$

For $Q_6^*$

Fix $p_6^*(0)$ and $P_6^*$ Let

$$H_6(z_6) = 1 - h_6(z_6)\int_0^\infty \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6} da_{,6}$$

And assume that

$$[k_3]_6 \int_0^\infty \beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(g_6^{-6}(b_6))a_6 - [k_2]_6 a_6}\, da_6 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_6(z_6)] > 0.$$

Clearly,

$$\frac{dH_6}{dz_6} = h_6(z)([f_2^{'}]_6(z_6))\int_0^{\infty} a_6\beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - f_6(P_6^*)a_6 - f_2(z_6)a_6}\, da_6$$

$-h^{'}(z_6)\int_0^{\infty}\beta_6(a_6)\, e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6}\quad da_6$

By the assumptions on $\left[f_2^6\right]_6 and\ [h^6]^6 dH_6/dz > 0$. Hence, if

$$H_6(0) = 1 - \int_0^{\infty}\beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6}da_6 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_6(0) < 0$ follows if $g_6^{-}(b_6) < [f_6^{-6}]_6(r_6)$, where $r_6$ is the solution of (12).

Indeed, since

$$R_6(s_6) = \int_0^{\infty}\beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 e^{-y}6 a}da_6$$

Is a decreasing function of $s_6$ when $P_6^* = [g^{-6}]_6(b_6) < [f_6^{-6}]_6(r_6)$, that is $[f_6]_6(P_6^*) < r$, $R\big(f_6(P^*)\big) > 1$which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_6(g^{-6}(b))) > 6, i.e.,$

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6\left(g_6^{-6}(b)\right)a_6} > 1. \quad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_6^*(a_6), Q_6^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_6^*(a_6), Q_6^*)$ is more complicated. We again consider perturbations $m_6(t_6 a_6) = p_6(t_6, a_6) - p_6^*(a_6)$ and $W_6(t) = Q_6(t) - Q^*$.Then

$$M_6(t) \equiv \int_0^\infty m_6(t, a_6) da_6 = P_6(t) - P_6^*,$$

$$[m_6]_6 + [m_a]_6 = [p_6]_6 + [p_a]_6 - \left[p^*\right]_6(a_6) \approx -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6$$

$$-[f_{6*}^6]_6 p_6^*(a_6)M - [f_{2*}^6]_6 p_6^*(\mathrm{a})W_6,$$

$$m_6(t,0)p_6(t,0) - p_6^*(0) \approx h_* \int_0^\infty \beta_6(a_6) m_6(t, a_6)\, da + \frac{[h_*^6]_6}{[h_*]_6}\, p_6^*(0) W_6$$

$$\frac{dW_6}{dt} = \frac{dQ_6}{dt} \approx [g_*^6]_6 Q_6^* M_{,6}$$

Where $[f_{6*}]_6$ denoting $[f_6(P^*)]_6$, $[f_{6*}]_6$ denoting $[f_2(P^*)]_6$, $\left[f_{6*}^{'}\right]_6$ denoting $\left[f^{'}{}_6(P^*)\right]_6 \left[f_{2*}^{'}\right]_6$ denoting $[f_2]_6$ $(Q_6^*,$ $h_{*}^{}{}_6$ denoting $h_*^{'}$ denoting $h_6^{'}$ $(Q_6^*)$, and $\left[g_*^{'}\right]_6$

Denoting $\left[g^{'}\right]_6(P_6^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_6]_6 + [m_a]_6 = -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6 - [f'_{6*}]_6[p_6^*](a_6)M_6 - [f'_{2*}]_6 p_6^*(a_6)W_6$$

$$m_6(t,0) = [h_*]_6 \int_0^\infty \beta_6(a_6)m_6(t,a_6)da_6 + \frac{[h_*^6]_6}{[h^*]_6}[p^*]_6(0)\,W_{,6} \qquad (30)$$

$$\frac{dW_6}{dt} = [g_*^6]_6[Q_6^*]M_6$$

Integrating $[m_6]_6(t,a_6)$ with respect to a from 0 to ∞, we have

$$\frac{dM_6(t)}{dt} = \int_0^\infty [mt]_6(t,a_6)da_6$$

$$= -\int_0^\infty [m_a]_6\,(t,a_6)da_6 - \int_0^\infty \mu(a_6)m(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + [f'_{6*}]_6 \int_0^\infty p_6^*(a_6)da_6\right)M_6 - [f_{2*}^6]_6 \int_0^\infty p_6^*(a_6)da_6 W_6$$

$$= m_6(t,0) - \int_0^\infty \mu_6(a_6)\, m_6(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + [f_{6*}']_6 P_6^*\right)M_6 - [f_{2*}']_6 P_6^* W_{,6}$$

From (30), we arrive at the equations

$$\frac{dW_6(t)}{dt} = \int_0^\infty ([h_*]_6\beta_6(a_6) - \mu_6(a_6))m_6(t,a_6)da_6$$

$$-\left([f_{6*}] + [f_{2*}]_6 + [f_{6*}']_6 P_6^*\right)M_6 + \left(\frac{[h_*']_6}{[h_*]^6}\,P_6^*(0) - f_{2*}' P_6^*\right)W_6,$$

(36)

$$\frac{dW_6(t)}{dt}\left[g^{'}{}_{*}\right]_6 Q_6^{*} M_{.6}$$

By the mean value theorem,

$$\int_0^{\infty}([h_{*}]_6\beta_6(a_6)-\mu_6(a_6))m_6(t,a_6)da_6=([h_{*}]_6\beta_6(\theta_6)-\mu_6(\theta_6))M_{,6}$$

Where $0<\ \theta_6=\theta_6(t)<\infty$.

And if $a_6^2\,[a_6^2]_6<4\,[a_2]_6<4\,a_2 e$,

$\mathrm{Re}\,\lambda_6\left[\lambda_6\right]_6=\mathrm{Re}\lambda_2\left[\lambda_2\right]_6=-\frac{6}{2}a_6[a2\ ]_6<0.$

Let $\Delta=\max_{2<u<\infty}\left|h_{*}\,[h_{*}]_6\beta_6(a_6)-\mu_6(a_6)\right|$. Then$|n_6(t)|$ 2Δ∀t 0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^{*}(a),\boldsymbol{Q}^{*})$ is asymptotically stable if either $a_6^2<4\,a_2 e$ and $[a_6]>2\Delta$,or $[a_6^2]$ 4$[a_2]$e and $(a_6-\sqrt{(a_6^2)_6-4(a_2)_6 e})>2\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a)=\mu_0+\left(\frac{a}{a+6}\right)$ and $\beta_6(a_6)=ae^{-ma}$ Then $\Delta=[\mu_0]+1$ Simple calculation leads to

$$\int_0^{\infty}\beta_6(a_6)e^{-\int_0^a\mu_6(\tau_6)+d\tau_6}da=\frac{6}{(m+\mu_0+1)^2}+\frac{2}{(m+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_6 + [\mu_0]$ 0.5. Assume$[f_6](x_6)= K_6x_6$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_6(x_6)= [0x_6]/(1 + x_6)$; the functional responses $[[f_2](x) = [0]\, x_6/(1 + x_6)$and $h_6(x_6)$=(1+$\omega_6 x_6$)/$(1 + x_6)$(See Holling [7] and Rashevsky [13])

Equation (64),now has the form of

$$\frac{m_6 + [\mu_0] + [\mathrm{K}][P_0])}{[m_6]_{\,6} + (\mu_0)_{\,6} + [\mathrm{K}][P_0])\,{}^{3}_{6}\ [\mathrm{K}]_{\,6}[P_0]\,{}^{3}_{6}} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(6 + P_0) < b_6$i e $[P_0] < b/(\delta_6 - b_6)$then the equilibrium$(P_0)$ $[P_0]_{\,6}(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_6$

$[f_{6*}] = K_6\ \mathrm{P}^{*}\ [f_{2*}] = \frac{\sigma \mathrm{Q}_6^*}{6+\mathrm{Q}_6^*}$ $[h_*] = \frac{6+\omega_6 \mathrm{Q}_6^*}{6+\mathrm{Q}_6^*}$

$[f_{6*}] = K_6\ \left[\mathrm{f}\,{}^{6}_{2*}\right] = \frac{\sigma_6}{(6+\mathrm{Q}_6^*)^2}[\ \mathrm{g}\,{}^{6}_{*}] = \frac{\sigma_6}{(6+\mathrm{P}_6^*)^2}$ $\left[\mathrm{h}\,{}^{6}_{*}\right] = \frac{\omega_6 - 6}{(6+\mathrm{Q}_6^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0,$ K, b, δ, and $\omega$ need to satisfy

$$\frac{\omega_6}{(m + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 + b_6} + \sigma)^{\,2}} + \frac{2\omega_6}{([m_6]_{\,6} + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 + b_6} + \sigma_6)^{\,3}} < 1$$

And

$$\frac{1}{(m_6 + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 - b_6})^{\,2}} + \frac{2\omega_6}{(m_6 + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 - b_6})^{\,3}} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_6 + [\mu_0]_6 + 1 + \frac{K_6 b_6}{\delta_6 - b_6} + (m_6 + [\mu_0]_6 + \frac{K_6 b_6}{\delta_6 - b_6} + \sigma_6)\, Q_6^*)^3$$

$$= m_6 + [\mu_0]_6 + 3 + \frac{K_6 b_6}{\delta_6 - b_6}(1 + Q^*)(1 + \omega_6 Q_6^*$$

Now

$$\left[a_6^2\right] - 4\left[a_2\right] e = \left(\frac{2K_6 b_6}{\delta_6 - b_6}\right)^2 + \frac{4K_6 b_6 \sigma_6 Q_6^*}{(\delta_6 - b_6)(1 + Q_6^*)} + \frac{\sigma_6^2 Q_6^{*2}}{(1 + Q_6^*)^2}$$

$$\frac{4\, b_6 \sigma_6 (\delta_6 - b) Q^*}{\delta_6 (1 + Q_6^*)^2} - \frac{4 (\delta_6 - b_6)^2 (1 - \omega_6) Q_6^*}{\delta_6 (1 + \omega_6 Q_6^*) 1 + Q_6^*)} P * (0)$$

And

$$\Delta = \lim_{0 < a < \infty} |[h_*]_6\, a e_6^{mu}\ [\mu_0]_6 - \frac{a_6}{a_6 + 6}$$

Hence if $[a_6^2 > 4(a_2)_6 e$ and

$$[a_6] = \frac{2K_6 b_6}{\delta_6 - b_6} + \frac{\sigma_6 Q_6^*}{6 + Q_6^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_6]_6 = [f_{6*}]_6 + [f_{2*}]_6 + [f_{6*}]\,P^*_6$$

$$[a_2]_6 = \left[f^6_{2*}\right]P^*_6 - \frac{h^6_*}{h_*}P^*(0)$$

$$e=[g^6_*]_6\,Q^*_6$$

$$\eta_6(t) = \left[h_*\right]_6\beta_6\left(\theta_6(t)\right) - \mu_6\left(\theta(t)\right),$$

The system(36)can be written as the following system:

$$[M(t)]\left[a_6\right]M_6 - [a_2]_6 W_6 - [n(t)]M_6$$

$$w_6(t) = eM_6$$

Write

$$A=\begin{pmatrix} -[a_6]_6 - [a_2]_6 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_6 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_6 \\ W_6 \end{pmatrix}$$

The solution of (32)is

$$X_6(t) = e^{A,t}\,[X]_6(0) + \int_0 e^{At}$$

Other problems to which the model is applied are: To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case β→∞ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness condition of (e) entropy. It is shown that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology

(gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure authors have discovered. It is shown that Legendre invariant thermodynamic scalar curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws, leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions,**

**geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged AdS black holes by considering the case of the RN-AdS black holes surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase transition. For the case $\omega_q=-2/3$ authors derive analytic expressions of (e) critical physical quantities, while (e) for cases $\omega_q\neq-2/3$ appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark

energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=)

chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross

section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions, while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center

collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma? Paul M. Chesler STOP Paul M. Chesler, Krishna Rajagopal** calculate how the energy and (e&eb) the opening angle of jets in N=4 SYM theory evolve as they propagate through (e&eb) the strongly coupled plasma of that theory. **Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb)

heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version) **Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb) strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion

collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar (=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI: 10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills

(SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the

decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed

black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to

equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a

nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about

the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic

map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstates thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11 August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical

advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.**

Characterizing the behaviour of strongly coupled quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at $\lambda=\lambda c$, the correlation length diverges as $\xi\sim|\lambda-\lambda c|-\nu$, while (e) the gap between the ground state and first excited state

vanishes as (e) ΔE∼ξ−z∼|λ−λc|zν. Exactly at the critical point λ=λc, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as ω∼|k−k0|z. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004 STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting

qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element. The basic architecture presented here can be expanded, offering (eb) the possibility for the coherent interaction of (e&eb) a large number of superconducting qubits. **Nature 449, 438-442 (27 September 2007) | doi:10.1038/nature06124; Received 18 April 2007; Accepted 25 July 2007 Coherent quantum state storage and transfer between two phase qubits via a resonant cavity Mika A. Sillanpää1, Jae I. Park1 & Raymond W. Simmonds1 STOP. Chui-Ping Yang, Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e)

long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007**

# PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION ATTRITION MODELS

## MOODULE NUMBER SEVEN

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Systems to which have applied the model are as follows: Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures have been proposed. An attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in**

**(eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP. In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011 http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 Qubit transfer**

between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys. Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer. Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP. In recent years S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P have investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys.

**Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP. Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2 [quant-ph] for this version) Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP. Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue from over (e&eb) a thousand plant and animal samples. NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP. Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this behavior, enhancing and accelerating (eb, eb+) the growth of correlations. Michael Kastner gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. Michael Kastner tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP. A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E.**

**Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91, 075435 – Published 26 February 2015 STOP. Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1 proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point'' with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. Membranes at quantum criticality Petr Hořava1, 2 Published 4**

**March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that (eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035 Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013 STOP Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study**

**the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen STOP Authors study the dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing**

**seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) Dynamic Service Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by H∗(X,(Λ∗TX)⊗g), where X is the target hyperk\"ahler manifold. We further give a mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the**

**low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition. Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb) the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537 © 2016 American Physical Society Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant (Δα/α) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and Δα/α to (e&eb) this conclusion.**

**However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts z≲3, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the Λ cold dark matter model value (jΛCDM=1); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in α(z), linked to (e&eb) a linear coupling (ζ) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order |ζ|∼10−4, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in Δα/α (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems, authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them**

**are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP Authors measure the star formation properties of (e) two large samples of galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the star formation rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a function of environment. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]. Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size √ n × √ n? The problem is that √n time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time O (√n) for d ≥ 3, or O (√nlog5/2n) for d = 2. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ∗This work was mostly done while the author was a PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship and by CIAR. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In**

**d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e) sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP The relationship between the foundations of mathematics and (e&eb) physics is a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb)**

**flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph] (or arXiv: 1503.05600v1 [math-ph] for this version) Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative**

**real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to the fine structure constant. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor goes to 0 as the (e&eb) time goes to 0, the big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph]**

**(or arXiv: 1306.4613v1 [quant-ph] for this version) Effects of gauge theory based number scaling on geometry Paul Benioff STOP In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x, outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph] (or arXiv: 1211.3381v2 [quant-ph] for this version) Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value, \int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances, the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an**

**observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include (e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP This work is based on (e) the field of reference frames based on (e) quantum representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul Benioff STOP A quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of**

**these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference: Journal of Physics: Conference Series 70 Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows: System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The**

**Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol**

## 41. MODEL FORMULATION

Let $p_7(t, a_7)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(g_7)Q_7$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_7(0) = 0, \quad g_7(P_7) > 0, \ P_7 \ 0, \qquad \lim_{P_7 \to \infty} g_7(P_7) = (K_7)_7 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_7 + (\rho_a)_7 + D_7(a_7, P_7, Q_7)\rho_7 + 0, \qquad 0 < a_7 \ \infty,$$

$$\rho_7(t,0) = \int_0^\infty n_7(a_7, P_7, Q_7)\rho_7(t_7, a_7)da_7$$

Where $D_7$the mortality or debit type function is function and $n_7$ is the fertility function.

Put $_7$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_7}{dt}\right]_7 = x_7 \left[\frac{dQ_7}{dt}\right]_7 - x_7 \left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_7}{dt}\right]_7 = x_3\left[\frac{dQ_3}{dt}\right]_7 - x_4\left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_7}{dt}\right]_7 - x_6\left[\frac{dQ_3}{dt}\right]_7$$

Where $x_7, x_7, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_7$ and $g_7(P_7)Q_7$

Now when you add the term $(b_7Q_7)$ with $g_7(P_7)(Q_7)$ we are infect see is that $(b_7Q_7)$ is credited incrementally part of $g_7(P_7)(Q_7)$ or instant nearly, when subtracted from $\frac{dQ_7}{dt}$ , it means that $g_7(P_7)(Q_7)$ is slowly eaten by $\frac{dQ_7}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

Following equation throws light on important issues in the following cases

**Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA**

(2509) (Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor.

(2510) In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface.

(2511) As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone.

(2512) Such planets are all the "right" distance from (e) their star, and hopefully, also have water.

(2513) But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone.

(2514) For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement.

(2515) But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway.

(2516) But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain.

(2517) It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life.

(2518) Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface.

(2519) The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113

**Abstract**

(2520) The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question.

(2521) Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry.

(2522) This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry.

(2523) Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments.

(2524) Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy.

(2525) The electronic and (e&eb) structural variability among them is extraordinary.

(2526) The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations.

(2527) The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface.

(2528) The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers.

(2529) Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed

mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp

**Water waves travel back in time to retrace their ripples By Jacob Aron (New Scientist)**

(2530) Nudge your bath tub just right and (eb) you might end up reversing time.

(2531) Don't worry – you won't end up wet, naked and without a towel in the Cretaceous period. Emmanuel Fort at the Langevin Institute in Paris, France, and his colleagues have developed what they call a "time mirror", which can make water waves rewind into (e&eb) their past. Humans and rubber duckies are (=) affected.

(2532) To understand how this works, imagine dropping a stone in a pond and watching waves ripple out from a central point in an ever-expanding circle. The equations governing the behaviour of waves say that it should also be possible for a ring of ripples to refocus back into a single point, like a reversed movie of the dropped stone. But we never actually see this happening, because it is vanishingly unlikely that an outside-in ring would form spontaneously. Put it in reverse

(2533) Fort and others have previously shown that it is possible to engineer the situation so the waves turn back on themselves by kitting out the edge of your pond with sensors and wave generators, effectively recording and then playing back the waves in reverse. The process has been done with water, sound and even microwaves, and is used to produce clearer ultrasound images for medical uses.

(2534) But now the team has discovered that simply dropping the tub downwards is enough to send water waves scurrying in reverse, with no need for complex sensors. That's because the velocity of water waves depends on acceleration due to gravity. "If you do a sudden jolt, you change the effective gravity, and the velocity of the wave changes," says Fort.

(2535) The disruption essentially resets the starting point of the wave. Once normal gravity is restored, it starts to travel both backwards and forwards. "Suddenly, you freeze the wave. When you release it again, it doesn't remember that it was going one way," says Fort.

(2536) Ripples in the water returned exactly to their start point - the shape of the Eiffel Tower

(2537) Ripples in the water returned exactly to their start point – the shape of the Eiffel Tower

**Déjà vu**

(2538) To show that the backward-travelling waves converge exactly on their starting point, just as if you were rewinding time, the team conducted a series of experiments. In some, they simply created a circular ripple of waves by dipping a probe in the water, similar to dropping a stone in the pond. In others, they used wave emitters in the shape of France, the Eiffel Tower and a smiley face to create complex wave-fronts. In each case, the reversed waves converged to recreate their original shape.

(2539) The experiments aren't truly reversing time, because the waves continue to travel out again once they have converged at the central point, points out Fort. "A real time-reversal would have a sink that absorbs all the waves," he says.

(2540) Now the team hopes to conduct similar experiments with other kinds of waves, like light or sound, although it may prove more difficult to find an analogue of their time-mirror jolt in other wave mediums, he says.

(2541) "The work is spectacular," says Geoffrey Edelmann of the US Naval Research Laboratory in Washington, D.C. "They essentially change the rules that govern how the wave propagates and, in doing so, create an entirely new way to focus waves free of spatial constraints." He also thinks the technique could be extended to other kinds of waves, and perhaps have medical applications such as lithotripsy, which uses ultrasound waves to break up kidney stones. Journal reference: Nature Physics, DOI: 10.1038/nphys3810

**Time reversal and holography with spacetime transformations Vincent Bacot, Matthieu Labousse, Antonin Eddi, Mathias Fink & Emmanuel Fort Nature Physics (2016) doi: 10.1038/nphys3810**

(2542) Wave control is usually performed by (e) spatially engineering the properties of a medium.

(2543) Because time and space play similar roles (e&eb) in wave propagation, manipulating time boundaries provides (eb) a complementary approach.

(2544) Here, authors experimentally demonstrate the relevance of this concept by (e) introducing instantaneous time mirrors.

(2545) Authors show with water waves that a sudden change of (e) the effective gravity generates (eb) time-reversed waves that refocus at the source.

(2546) They generalize this concept for (e) all kinds of waves, introducing (eb) a universal framework which explains (eb) the effect of any time disruption on (e&eb) wave propagation.

(2547) Authors show that sudden changes of the medium properties generate (eb) instant wave sources that emerge (eb) instantaneously from the entire space at (eb) the time disruption.

(2548) The time-reversed waves originate from (e) these 'Cauchy sources', which are (+) the counterpart of Huygens virtual sources on (eb) a time boundary. It allows us to revisit the holographic method and introduce a new approach for wave control.

**Viewpoint: A Bird's Eye View of Circuit Photons Sebastian Schmidt, Institute for Theoretical Physics, ETH Zurich, 8093 Zurich, Switzerland June 30, 2016• Physics 9, 74 A scanning probe detects the quantum states of photons in a microwave circuit, providing the information needed for quantum simulations. Figure caption expand figure A. Houck/Princeton (deleted)**

(2549) Figure 1: Scanning defect microscopy provides a map of photons in a resonator lattice. Houck and colleagues demonstrated the technique using (e) 49 resonators (grey lines) that were coupled together to (e&eb) form a kagome lattice.

(2550) Quantum mechanics rules the dynamics of (e) light and matter.

(2551) Yet performing a quantum-mechanical simulation of a material from first principles is (=) practically impossible on a classical computer because (e) the complexity of the simulation increases (eb) exponentially with the number of particles involved.

(2552) The solution, according to Richard Feynman, was to build a machine out of (e) quantum building blocks that could directly emulate (e&eb) the material itself [1].

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| (2553) Prototypes of such quantum simulators that are based on (e) ultracold atoms, ions, photons, and superconducting microwave circuits are now available [2], with (e&eb) the latter, in particular, having attracted (e) Silicon Valley's interest. |
| (2554) The challenge with these circuit-based simulators, however, is that they are (=) 2D, which complicates the readout of their constituent elements. Andrew Houck from Princeton University, New Jersey, and colleagues have now delivered an attractive solution by developing a technique [3], called scanning defect microscopy (Fig. 1), that determines (eb) the number of photons occupying (e) each mode of a 2D microwave circuit. |
| (2555) It is this information that would serve as (=) the fundamental input and output for (e) certain quantum simulations. |
| (2556) Superconducting microwave circuits combine electronic and (e&eb) photonic degrees of freedom [4, 5]. |
| (2557) The main element of the circuit is (=) a transmission line, which is made up of (e) a central superconducting wire separated by a gap from (e) two grounded plates. |
| (2558) All of these structures are on a single plane, as if one had taken a 2D slice through (e&eb) a coaxial cable. |
| (2559) When truncated, the transmission line becomes (=) a resonator, which can host (e&eb) discrete photon modes within (eb) its gaps |
| (2560) Large lattices of resonators can be engineered in various 1D or 2D geometries by coupling (e&eb) two, three, or more resonators together via (e&eb) a capacitive interface. In many ways, photons in such devices behave similarly to electrons in a solid. |
| (2561) To make microwave circuits that can simulate (e&eb) quantum phenomena faster than a classical computer, however, resonator lattices have to be integrated with (e&eb) superconducting qubits |
| (2562) . Such qubits are controlled with (e&eb) electrical currents in a Josephson tunnel junction, and in many respects, they behave (e&eb) like artificial atoms, which couple to (e&eb) the photons in the circuit. |
| (2563) As a result, superconducting microwave circuits can be used to (e) explore the coupling between the quantum states of light and (e&eb) matter, the regime of (e) circuit quantum electrodynamics (cQED). |
| (2564) Photons in these devices often exhibit (eb) striking matter-like behavior [6, 7], providing (eb) the basis for the simulation of complex materials. |
| (2565) Such circuits can be fabricated on a substrate using (e) standard lithographic techniques, with (e&eb) qubits and resonators that are (=) hundreds of micrometers or even millimeters in size. |
| (2566) Small-scale realizations of cQED lattices have already been used for (e) proof-of-principle quantum simulations [8, 9]. |
| (2567) However, future large-scale simulations will have to be initialized and measured (eb) efficiently—steps known as (=) drive and readout. |
| (2568) This requires connecting the lattices to (e) external circuitry, such as voltage sources, circulators, and amplifiers, which typically reside outside (e) the chip. |
| (2569) So far, this connectivity has been realized through (e&eb) external transmission lines, but these elements can only couple to (e&eb) the resonators at the edges of (e) the lattice because of (e) its 2D nature. |
| (2570) Such a side view provides (eb) only limited information about (e&eb) the quantum state and spatial correlations of photons inside (eb) the bulk resonators, seriously limiting (e) the applicability of circuit QED devices for (e) quantum simulations. |

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| (2571) The probe developed by Houck and colleagues, scanning defect microscopy (SDM), solves (e) this problem by (e) moving outside the plane of the chip to provide (eb) a fuller image of photons within (eb) the chip. |
| (2572) In SDM, a small sapphire tip is positioned a few tens of (e) micrometers above the surface of one of the resonators (Fig. 1). |
| (2573) The presence of this dielectric modifies (e&eb) the evanescent electric field from the circuit photons, which extends (eb) just outside the plane of the chip. |
| (2574) Specifically, the tip increases (eb) the local capacitance of the transmission line, downshifting (e) the local resonator frequency. |
| (2575) This in turn causes (eb) a drop in the transmission of photons through (e&eb) the lattice when (e) it is driven at its nominal resonance frequency. |
| (2576) Houck's team showed that if the sapphire tip has only a weak (linear) effect on the resonator frequency, then (eb) a simple analytic relationship exists between the induced frequency shift, (e&eb) the change in transmission, and (e&eb) the local photonic occupation of each resonator. |
| (2577) This relationship provides (eb) a recipe for detecting the local photon occupancy inside (eb) the lattice. |
| (2578) First, calibrate the probe by determining (eb) the frequency shift for (e) a single isolated resonator as (=) a function of the tip's position. |
| (2579) Second, measure the transmission through (e&eb) the whole lattice with and without (e&eb) the probe. |
| (2580) Finally, use the measured transmission drop with (e&eb) the probe present to find (eb) the photonic occupation of (e) an individual resonator in (eb) the lattice. |
| (2581) A map of the photon-mode occupancies can be obtained by (e) scanning the probe across the whole lattice. |
| (2582) Houck's team used (e) SDM to image the occupancies of (e) the photon modes of (e) 49 coupled resonators arranged in (eb) a special geometry called a kagome lattice (see details in Fig. 1). |
| (2583) (To simplify the measurement, no superconducting qubits were (=) incorporated with (e&eb) the resonators.) |
| (2584) The occupancies they measured were in excellent agreement with (=) theoretical predictions, corresponding to (e&eb) a fidelity of 0.99. |
| (2585) The high fidelity is even more impressive considering the weak signal-to-noise ratio of (e) their technique (which is limited by (e) the intentionally small size of the frequency-shifted region) and (e&eb) the technical challenge of (e) positioning the probe inside (eb) a cryostat for (e) in situ measurements. |
| (2586) While SDM shares similarities with (e&eb) so-called bead-pull experiments for accelerators and scanning gate microscopy for (e) nanostructures, it is truly a (=) new tool for quantum electronics. |
| (2587) First application of SDM is likely to be (=) the simulation of 2D topological phenomena and (e&eb) geometrically frustrated lattices (of which the kagome lattice is (=) one example.) |
| (2588) Both are associated with (e&eb) macroscopically degenerate systems, which are (=) notoriously difficult to simulate on (e&eb) a classical computer. |
| (2589) Such circuits are (=) genuinely open systems (photons can be added or lost). |
| (2590) This combined with (e&eb) their inherent strong light-matter coupling, makes (eb) such quantum simulators ideal candidates for (e) the study of many-body, out-of-equilibrium phenomena, such as (=) thermalization or nonequilibrium phase transitions. |

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| (2591) SDM could also lead to (eb) the discovery of new exotic phenomena, which are (=) unknown in traditional solid-state systems. |
| (2592) In a nutshell, SDM constitutes a new and indispensible tool for (e) future experiments on quantum simulations with microwave photons. This research is published in Physical Review X. |

**Imaging Photon Lattice States by Scanning Defect Microscopy D. L. Underwood, W. E. Shanks, Andy C. Y. Li, Lamia Ateshian, Jens Koch, and A. A. Houck Phys. Rev. X 6, 021044 – Published 30 June 2016**

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| (2593) Microwave photons inside (eb) lattices of coupled resonators and superconducting qubits can exhibit (eb) surprising matter like behavior. |
| (2594) Realizing such open-system quantum simulators presents (eb) an experimental challenge and requires (e) new tools and measurement techniques. |
| (2595) Here, authors introduce scanning defect microscopy as one such tool and illustrate its use in mapping (e&eb) the normal-mode structure of microwave photons inside (eb) a 49-site kagome lattice of coplanar waveguide resonators |
| (2596) Scanning is accomplished by (e) moving a probe equipped with a sapphire tip across (e&eb) the lattice. |
| (2597) This locally perturbs (e&eb) resonator frequencies and induces (eb) shifts of the lattice resonance frequencies, which authors determine by (e) measuring the transmission spectrum. |
| (2598) From the magnitude of mode shifts, author can reconstruct (eb) photon field amplitudes at (eb) each lattice site and thus create (eb) spatial images of the photon-lattice normal modes. Received 4 December 2015 DOI:http://dx.doi.org/10.1103/PhysRevX.6.021044 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society |

| |
|---|
| |
| We assume that |
| $$D_7 = \mu_7(a_7) + (f_7)_7(P_7)_7 + (f_7)_7(Q)_7 \quad (7)$$ |
| $$n_7 = \beta_7(a_7)h_7(Q_7), \quad (3)$$ |
| Where $\mu_7$the age- dependent debit rate whit is $\mu_7(a_7)$ 0 $\forall a > 0$; $(f_7)_7$ is the density- dependent death rate with $f_7(0) = 0$, $(f_7^7)_7(P_7)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_7)_7(P_7) = (k_7)_7$ ; $f_7$ is the functional response of creditors to prey with $[f_7]_7(0) = 0, (f_7^7)_7(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_7)_7(Q_7) = [k_7]$ ; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^7(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0. |

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_7 + (\rho_a)_7 = -\mu_7(a_7) + (f_7)_7 P_7) + \big(f_7)_7(Q_7)\big)(\rho_7)_7 \quad (4)$$

$$\rho_7(t_7, 0) = h_7(Q_7) \int_0^\infty \beta_7 (a_7)\rho_7(t, a_7)da, \quad (5)$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(P_7)Q_7 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_7$ and $Q_7$ of (4-6) satisfy

$$0 \leq \rho_7(t, a_7) \leq \rho_7^*(t, a_7),$$

$$0 \leq Q_7(t) \leq Q_7^*(t),$$

Where $\rho_7^*(t, a_7)$ and $Q_7^*(t)$, are solutions of

$$(\rho_t^*)_7 + (\rho_a^*)_7 = -(\mu_7(a_7) + (f_7)_7(P_7^*) + (f_7)_7 Q_7^*))\, \rho_7^*, \quad (7)$$

$$\rho_7^*(t, 0) = \int_0^\infty \beta_7(a_7)\, \rho_7^*(t, a_7)da_7 \quad (8)$$

$$\frac{dQ_7^*}{dt} = -b_7\, Q_7^* + g_7(P_7^*)Q_7^*, \quad (9)$$

With $P_7^*(t) = \int_0^\infty \rho_7^*(t, a_7)da_7$

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From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_7^*(t)}{dt} = P_7^*(t)\big(r_7 - (f_7)(P_7^*) - (f_7)_7(Q^*)\big), \quad (10)$$

$$\frac{dQ_7^*(t)}{dt} = Q_7^*(t)(-b_7 + g_7(P_7^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_7(a)e^{-\int_0^a \mu_7(\tau)dt}e^{-ra}\,da_7 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_7^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(t)dt}\,da_7 > 1, \quad (H1)$$

Which guarantees $r_7 > 0$.

*Remark.* If $\beta_7(a) = [\beta_0]_7 a_7 e^{-\alpha_7 a_7}, \mu_7(a_7) \equiv \mu_7$ then the solution of (12) is

$$r_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_7]_7(P_7) = 0, [f_7](Q_7) = \lambda_7 Q_7\ h_7(Q_7) = 1, \quad and\ \ g_7(P_7) = c_7 P_7$$

$$[\rho_7]_7 + [\rho_a]_7 = -(\mu_7 + \lambda_7 Q_7)\rho_7$$

$$\rho_7(t,0) = \int_0^\infty [\beta_0]_7 a_7 e^{-\alpha_7 a_7} \rho_7(t,a_7)\, da_7$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + c_7 P_7 Q_7$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_7}{dt} = \omega_7 P_7 - \lambda_7 P_7 Q_7$$

$$\frac{dQ_7}{dt} = b_7 Q_7 - c_7 P_7 Q_7$$

Where $\omega_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$

## 42. THE EQUILIBRIUM $(\rho_0)_7(a_7), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_7(a_7), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_7 + (\rho_a)_7 = -\big(\mu_7(a_7) + (f_7)_7(P_7)\big)\rho,$$

$$\rho_7(t,0)$$

$$= \int_0^\infty \beta_7(a_7)\rho_7(t,a_7)da_7 \qquad (13)$$

The solution of (13) has the form of

$$\rho_7(t,a_7) = Ce_7^{r(t-a_7)} - \int_0^a (\tau)d\tau - \int_0^7 (f_7)_7\big(P_7(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_7(t) = C_7 e^{r_7 t - \int_0^t (f_7)_7(P_7(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_7^{-7}(r)$ if $k_7 > r$. This is equivalent to

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{(k_7)_7 a_7} da_7 < 1.$$

In fact, if we define

$$R_7(s_7) = \int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau)d\tau} e^{-s_7 a_7}\, da_7$$

$R(s)$ is a decreasing function. $R_7(r_7) = 1$ And $R_7(k_7)_7 < 1$ imply $r_7 < (k_7)_7$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_7(a_7), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{k_7 a_7} da_7$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_7 \equiv \int_0^\infty (\rho_0)_{77}(a_7)da_7$ *satisfies*

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{-(f_7)_7(P_0)a_7} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_7(a_7), 0)$

Let $u_7(t, a_7) = \rho(t, a_7) - \rho_0(a_7)$ Then

$$N_7(t) \equiv \int_0^\infty u_7(t, a_7) da_7 = \int_0^\infty \rho_7(t, a_7) da_7 - \int_0^\infty (\rho_0)_7(a_7) da = P_7(t) - P_0)_7$$

And

$$(u_7)_7 + (u_a)_7 = (\rho_t)_7 + (\rho_a)_7 - (\rho_0^7)_7(a_7)$$

$$= -\mu_7(a_7) + (f_7)_7(P_7) + (f_7)_7(Q_7))(u_7 + (\rho_0)_7) + \mu_7(a_7)(\rho_0)_7 + (f_7)_7(P_0)_7(\rho_0)_7$$

$$= -\mu_7(a_7)u_7 - (f_7)_7(P_7) + (f_7)_7(Q_7))\, u_7 - (\rho_0)_7\, (f_7)_7(P_7) - (f_7)_7(P_0)_7$$

$$+ (f_7)_7(Q_7)), \qquad (15)$$

Where $(\rho_0^7)_7(a_7)$ denotes $d(\rho_0)_7(a_7)/da_7$

Expansion $(f_7)_7(P_7)$ and $(f_7)_7(Q_7)$ at $(P_0)_7$ and 0 yield respectively

$$(f_7)_7(P_7) = (f_7)_7((P_0)_7 + N_7)$$

$$= (f_7)_7(P_0)_7 + (f_7^7)_7(P_0)_7 N_7 + o_7(N_7), \qquad (16)$$

$$(f_7)_7(Q_7)$$

$$= (f_7^7)_7(0)Q + o_7(Q_7) \qquad (17)$$

Write $f_7(P_0) = f_{70}, f_7'(P_0) = f_{70}'$, and $f_7'(0) = f_{70}'$ and substitute (16) and (17) into (15) to obtain

$$(u_7)_7 + (u_a)_7 \approx -(\mu_7(a_7) + (f_{70})_7)u - (f_{70}^7)_7(\rho_0)_7(N_7) - (f_{70}^7)_7(\rho_0)_7(Q_7)$$

$$u_7(t, 0) = \rho_7(t, 0)(\rho_0)_7(0)$$

$$= h_7(Q_7)\int_0^{\infty} \beta_7(a_7)(u_7(t,a) + (\rho_0)_7(a_7)da_7 - (\rho_0)_7\,(0)$$

$$= h_7(Q_7)\int_0^{\infty} \beta_7(a_7)(u_7(t,a_7)da_7 + (h_7(Q_7) - 1)(\rho_0)_7\,(0) \qquad (18)$$

Expanding

$$h_7(Q_7) = 1 + \ (h^7)_7(0)Q_7 + 0(Q_7) \equiv 1 + h_0^{'}Q_7 + 0Q_7$$

And substituting it into (18), we arrive at

$$u_7(t,0) \approx \int_0^{\infty} \beta_7\,(a_7)u_7(t,a_7)da_7 + (h_{a_7}^{'})(\rho_0)_7Q_7$$





The linearization is the following :

$$[u_7] + [u_7] = -(\mu_7(a_7) + [t_{70}]_7u_7 - [\alpha_7]_7(a_7)N_7 - [\alpha_7]_7(a_7)Q_7$$

$$u_7(t,0) = \int_0^{\infty} \beta_7(a_7)u_7(t,a_7)da_7 + [\xi_7]Q_7 \qquad (19)$$

$$\frac{dQ_7}{dt} = -b_7Q_7$$

where

$$[\alpha_7]_7(a_7) = [t_7^7a]_7[p,o]_7(a_7)d2(a_7) = [f_{70}^7]_7(\rho\ o\ )_7(a_7)\xi = [\mathrm{h\ o}^7]_7[\rho\mathrm{o}]_7(\mathrm{o}) \text{ and } \mathrm{b}_7 = \mathrm{b}_7 - g_7(P_7)$$

If $g_7(P_0) > b$, clearly, $(\rho_0)_7(a_7), 0)$ is unstable and the predator population is persistent. Assume $g_7(P_0)_7 < b$.

Then, $Q_7(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_7 + u_a = -(\mu_7(a_7) + (f_{70})_7)u_7 - (\alpha_7)_7(a_7)N_7$$

$$= -(\mu_7(a_7) + f_{70})u_7 - \alpha_7(a_7)\int_0^\infty u_7(t, a_7)da_7$$

$$u_7(t, 0) = \int_0^\infty \beta_7\,(a_7)u_7(t, a_7)da_7$$

Using the method developed by Cushing [2, 3], we suppose

$$u_7(t, a_7) = \tilde{u}_7(a_7)^{e^{\gamma(t-a_7)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_7}\,da_7$. Then,

A solution of the first differential equation for $\tilde{u}(a_7)$ gives,

$$\tilde{u}_7(a_7) = e^{-\int_0^a \mu_7(\tau_7)d\tau - (f_{70})_7 a}\left(\tilde{u}_7(0) - w\int_0^a e^{\int_0^\tau \mu_7(s_7)ds + (f_{70})}\tau_7\,(\alpha_7)_7(\tau_7)e^{\gamma_7\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_7(a_7)e^{-\gamma_7 a_7 - (f_{70})_7 a_7 - \int_0^a \mu_7(\tau_7)d\tau}\right)\tilde{u}_7(0)$$

$$+w\int_0^\infty \beta_7(a_7)e^{-\gamma_7 a_7 - (f_{70})_7 a_7 - \int_0^a \mu_7(\tau_7)d\tau}\int_0^a (\alpha_7)_7(\tau_7)e^{-\gamma_7\tau_7}e^{(f_{70})_7\tau_7 + \int_0^\tau \mu_7(s_7)ds}d\tau_7\,da_7$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty} e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau} da\right)(\tilde{u}_7(0))_7$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau}\int_0^a e^{(f_{70})_7\tau_7+\int_0^{\tau}\mu_7(s_7)ds+y\tau}\alpha_7(\tau)d\tau\, da\right)w=0. \quad (23)$$

Define $Y(\gamma,a)=e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(t_7)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)da_7\right)u_7(0)$$

$$+w\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\, da_7=0, \quad (24)$$

$$\left(\int_0^{\infty}Y_7(\gamma_7,a_7)da_7\right)\tilde{u}_7(0)$$

$$+\left(-1-\int_0^{\infty}Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\, da_7\right)w=0.$$

Equation (20) has a solution of the form of $u_7(t,a_7)=u_7(a_7)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_7(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_7(\gamma_7,a_7)=\int_0^a(\rho_0)_7(\tau_7)Y^{-7}(\gamma_7,a_7)d\tau_7$. We obtain the characteristic equation

$$f_{70_7}^{l}\left(\int_0^{\infty} Y_7(\gamma_7,a_7)da_7\right)\left(\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$+\left(1-\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)da\right)\left(1+(f_{70}^{l})_7\int_0^{\infty}Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_7(a_7)=(\rho_0)_7(0)e^{-\int_0^a \mu_7(\tau_7)d\tau_7-(f_{70})_7a_7}$$

$$I_7(\gamma_7,a_7)=\int_0^a(\rho_0)_7(\tau_7)e^{\gamma_7\tau_7+(f_{70})_7\tau+\int_0^{\tau}\mu_7(s_7)ds_7}d\tau$$

$$=\frac{(\rho_0)_7(0)}{\gamma_7}(e^{\gamma_7a_7}-1)$$

Let

$$E_7(a_7)=e^{-(f_{70})_7a_7-\int_0^a\mu_7(\tau_7)d\tau_7}$$

Then $[Y_7]_7[\gamma_7]_7,a_7)=E_7(a_7)e^{-\gamma_7a_7}$. by a simple calculation,

$$\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da=\frac{[\rho_0]_7(0)}{\gamma}\left(1-\int_0^{\infty}\beta_7(a_7)E_7(a_7)e^{-\gamma_7a_7}\ da\right);$$

$$\int_0^{\infty}Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da=\frac{[\rho_0]_7(0)}{\gamma_7}\left(\int_0^{\infty}E_7(a_7)e^{-\gamma_7a_7}\ da_7\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_7(0)}{\gamma_7}[f_{70}^7]\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\,(\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7)$$

$$+\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right).$$

$$\left(1+\frac{(\rho_0)_7(0)}{\gamma_7}(f_{70}^{'})_7\left(\int_0^\infty E_7(a_7)da_7-\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right)(1+\frac{(f_{70}^{'})_7}{\gamma_7}(\rho_0)_7(0)\int_0^\infty E_7(a_7)da_7)=0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)da_7 = 1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7 \neq 0,$$

Since $\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)\int_0^a \alpha_7(\tau_7)Y^{-7}(\gamma_7,\tau_7)d\tau_7 da_7 > 0$. Therefore, the characteristic solution is

$$\gamma_7 = -(f_{70}^{'})_7(\rho_0)_7\int_0^\infty E_7(a_7)da_7.$$

Since $f_7'(P_0)_7 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_7), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE OF DEBIT AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_7), 0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_7 + (\mu_7(a_7) + [f_7]_7(P_7^*) + [f_7]_7(Q_7^*))P_7^* = 0, \quad (76)$$

$$p_7^*(0) = h_7(Q_7^*)\int_0^\infty \beta_7(a_7)p_7^*(a_J)p_7^*(a_7)da_7 \quad (77)$$

$$g_7(P_7^*) - b_7 = 0, \quad (78)$$

Where $P_7^* = \int_0^\infty p_7^*(a_7)\, da_7$

Equation (28) yields $P^* = g^{-7}(b)$,

From (26),

$$p_7^*(a) = p_7^*(0)e^{\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(Q_7^*)a_7}$$

Substituting into (27), we need to solve

$$\left(1 - h_7(Q_7^*)\int_0^{\infty} \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(Q_7^*)a_7} da_7\right) p_7^*(0) = 0$$

For $Q_7^*$

Fix $p_7^*(0)$ and $P_7^*$ Let

$$H_7(z_7) = 1 - h_7(z_7)\int_0^{\infty} \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(z_7)a_7} da_{,7}$$

And assume that

$$[k_3]_7 \int_0^{\infty} \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(g_7^{-7}(b_7))a_7 - [k_7]_7 a_7}\, da_7 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_7(z_7)] > 0.$$

Clearly,

$$\frac{dH_7}{dz_7} = h_7(z)([f_7']_7(z_7))\int_0^{\infty} a_7\beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - f_7(P_7^*)a_7 - f_7(z_7)a_7} da_{7}$$

$-h'(z_7)\int_0^{\infty} \beta_7(a_7)\, e^{\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(z_7)a_7}\ da_7$

(See author's It is a debit-credit world)

By the assumptions on $\left[f_7^7\right]_7 and\ [h^7]^7 dH_7/dz > 0$. Hence, if

$$H_7(0) = 1 - \int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7} da_7 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_7(0) < 0$ follows if $g_7^-(b_7) < [f_7^{-7}]_7(r_7)$, where $r_7$ is the solution of (12).

Indeed, since

$$R_7(s_7) = \int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 e^{-y\frac{a}{7}}} da_7$$

Is a decreasing function of $s_7$ when $P_7^* = [g^{-7}]_7(b_7) < [f_7^{-7}]_7(r_7)$, that is $[f_7]_7(P_7^*) < r$, $R\left(f_7(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

**THEOREM 3**

Assume that the hypothesis (H3) holds and $R(f_7(g^{-7}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7\left(g_7^{-7}(b)\right)a_7} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_7^*(a_7), Q_7^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_7^*(a_7), Q_7^*)$ is more complicated. We again consider perturbations $m_7(t_7 a_7) = p_7(t_7, a_7) - p_7^*(a_7)$ and $W_7(t) = Q_7(t) - Q^*$.Then

$M_7(t) \equiv \int_0^\infty m_7(t,a_7)da_7 = P_7(t) - P_7^*,$

$$[m_7]_7 + [m_a]_7 = [p_7]_7 + [p_a]_7 - [p^*]_7(a_7) \approx -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7$$

$$-[f_{7*}^7]_7 p_7^*(a_7)M - [f_{7*}^7]_7 p_7^*(\mathrm{a})W_7,$$

$$m_7(t,0)p_7(t,0) - p_7^*(0) \approx h_* \int_0^\infty \beta_7(a_7)m_7(t,a_7)\,da + \frac{[h_*^7]_7}{[h_*]_7}\, p_7^*(0)W_7$$

$$\frac{dW_7}{dt} = \frac{dQ_7}{dt} \approx [g_*^7]_7 Q_7^* M,_7$$

Where $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}']_7$ denoting $[f'_7(P^*)]_7[f_{7*}']_7$ denoting

$[f_7]_7$ $(Q_7^*,$ $h_{7}^*$ denoting $h_*'$ denoting $h_7'$ $(Q_7^*)$, and $[g_*']_7$

Denoting $[g']_7(P_7^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$[m_7]_7 + [m_a]_7 = -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7 - [f'_{7*}]_7[p_7^*](a_7)M_7 - [f_{7*}']_7 p_7^*(a_7)W_7$

$m_7(t,0) = [h_*]_7 \int_0^\infty \beta_7(a_7)m_7(t,a_7)da_7 + \frac{[h_*^7]_7}{[h^*]_7}[p^*]_7(0)\,W,_7$ (30)

$$\frac{dW_7}{dt} = [g_*^7]_7[Q_7^*]M_7$$

Integrating $[m_7]_7(t,a_7)$ with respect to a from 0 to ∞, we have

$$\frac{dM_7(t)}{dt} = \int_0^\infty [mt]_7(t, a_7)da_7$$

$$= -\int_0^\infty [m_a]_7\,(t, a_7)da_7 - \int_0^\infty \mu(a_7)m(t, a_7)da_7$$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f^{'}{}_{7*}]_7 \int_0^\infty p_7^*(a_7)da_7\right)M_7 - [f_{7*}^7]_7 \int_0^\infty p_7^*(a_7)da_7 W_7$$

$$= m_7(t, 0) - \int_0^\infty \mu_7(a_7)\, m_7(t, a_7)da_7$$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 - [f_{7*}^{'}]_7 P_7^* W_{.7}$$

From (30), we arrive at the equations

$$\frac{dW_7(t)}{dt} = \int_0^\infty ([h_*]_7\beta_7(a_7) - \mu_7(a_7))m_7(t, a_7)da_7$$

$$-\left([f_{7*}] + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 + \left(\frac{[h_*^{'}]_7}{[h_*]^7}\, P_7^*(0) - f_{7*}^{'} P_7^*\right)W_7,$$

(31)

$$\frac{dW_7(t)}{dt}[g^{'}{}_*]_7\, Q_7^*\, M_{.7}$$

By the mean value theorem,

$$\int_0^\infty ([h_{"*}]_7\beta_7(a_7) - \mu_7(a_7))m_7(t, a_7)da_7 = ([h_*]_7\beta_7(\theta_7) - \mu_7(\theta_7))M,_7$$

Where $0< \theta_7 = \theta_7(t) < \infty$.

And if a $_7^7$ [a $_7^7$]2<4 [a $_7$]2<4 a $_7$e,

Re $\lambda_7[\lambda_7]_7 = \mathrm{Re}\lambda_7[\lambda_7]_7 = -\frac{7}{7}a_7[\mathrm{a}2\ ]_7 < 0$.

Let ⊿=max$_{7<u<\infty}$ |h $_*$ [h $_*$] $_7$β $_7$($a_7$) − $\mu_7(a_7)$|. Then|$n_7$(t)| ⊿∀t . Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium(ρ $^*$ (a), $\boldsymbol{Q}^*$) is asymptotically stable if either a $_7^7$ <4 a $_7$ e and [a $_7$] > 1⊿,or [a $_7^7$] [a $_7$]e and (a $_7$ − $\sqrt{(\mathrm{a}_7^7)_7 - 4(\mathrm{a}_7)_7\mathrm{e})} > 1⊿$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{\mathrm{a}}{\mathrm{a}+7}\right)$ and β $_7$($a_7$) = ae-$^{\mathrm{ma}}$ Then ⊿=[$\mu_0$] + 2 Simple calculation leads to

$$\int_0^\infty \beta_7(a_7)\mathrm{e}^{-\int_0^{\mathrm{a}}\mu_7(\tau_7)+d\tau_7\mathrm{da}} = \frac{1}{(\mathrm{m}+\mu_0+1)^7} + \frac{1}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_7$ + [$\mu_0$] .5. Assume[$f_7$]($x_7$)= $K_7x_7$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_7(x_7)$= [0$x_7$]/(2 + $x_7$); the functional responses [[$f_7$]($x$) = [0] $x_7$/(1 + $x_7$)and $h_7(x_7)$=(1+$\omega_7x_7$)/(1 + $x_7$)(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_7 + [\mu_0] + [\mathrm{K}][P_0])}{[m_7]_7 + (\mu_0)_7 + [\mathrm{K}][P_0])_7^3\ [\mathrm{K}]_7[P_0]_7^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1+P_0) < b_7$ i e $[P_0] < b/(\delta_7 - b_7)$then the equilibrium$(P_0)$ $[P_0]_7(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_7$

$[f_{7*}] = K_7\ \mathrm{P}^*\ [f_{7*}] = \frac{\sigma \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*} \quad [h_*] = \frac{7+\omega_7 \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*}$

$[f_{7*}] = K_7\ [\mathrm{f}_{7*}^{\ 7}] = \frac{\sigma_7}{(7+\mathrm{Q}_7^*)^7} [\ \mathrm{g}_*^{\ 7}] = \frac{\sigma_7}{(7+\mathrm{P}_7^*)^7} \quad [\mathrm{h}_*^{\ 7}] = \frac{\omega_7 - 7}{(7+\mathrm{Q}_7^*)^7}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_7}{(m + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma)^7} + \frac{1\omega_7}{([m_7]_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma_7)^3} < 1$$

And

$$\frac{1}{(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^7} + \frac{1\omega_7}{(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $\mathrm{Q}^*$ which can be solved from the following equation:

$(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7} + (m_7 + [\mu_0]_7 + \frac{K_7 b_7}{\delta_7 - b_7} + \sigma_7)\ \mathrm{Q}_7^*)^3$

$= m_7 + [\mu_0]_7 + 3 + \frac{K_7 b_7}{\delta_7 - b_7}(1 + \mathrm{Q}^*)(1 + \omega_7 \mathrm{Q}_7^*$

Now

$$\left[\mathrm{a}\begin{matrix}1\\1\end{matrix}\right]-4\left[\mathrm{a}_{7}\right]\mathrm{e}=\left(\frac{1K_7 b_7}{\delta_7-b_7}\right)^{7}+\frac{4K_7 b_7\sigma_7\,\mathrm{Q}_{7}^{*}}{(\delta_7-b_7)(7+\mathrm{Q}_{7}^{*})}+\frac{\sigma_{7}^{7}\mathrm{Q}_{7}^{*7}}{\left(1+\mathrm{Q}_{7}^{*}\right)^{7}}$$

$$\frac{4\,b_7\sigma_7\,(\delta_{7}-\mathrm{b})\mathrm{Q}^{*}}{\delta_{7}\,(1+\mathrm{Q}_{7}^{*})^{7}}-\frac{4\,(\delta_{7}-b_7)^{7}\,(1-\omega_7)\mathrm{Q}_{7}^{*}}{\delta_{7}\left(1+\omega_7\mathrm{Q}_{7}^{*}\right)(1+\mathrm{Q}_{7}^{*})}\mathrm{P}^{*}(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_{7}\,ae_{7}^{mu}\;\;[\mu_0]_{7}-\frac{a_7}{a_7+7}$$

Hence if $[\mathrm{a}_{7}^{7}>4(a_7)_{7}\mathrm{e}$ and

In the following are given some more modular structures that are applied with the model and the concomitant results derived thereof:

Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments

(1) We have an intuitive understanding of how the world works. Our intuition works well in our everyday lives, but it doesn't quite match (e&eb) reality.

(2) One of the basic assumptions authors make about objects is that they are (=) objectively real.

(3) For example, when you leave your apartment in the morning, you imagine your house is right where you left it as you go about your day. Barring some calamity, you expect to find your apartment just as you left it when (e) you return in the evening.

(4) This idea is known as Macrorealism, in that a macroscopic object such as your house is still there even when (e) you're not looking at it

(5) In scientific terms we would say the state of an object exists independent of (e) an observer. While this seems obviously true, we can test the idea experimentally.

(6) To show this, let's consider an experiment. Suppose we have a box that contains a coin. At any time we can open the box and (eb) see whether it's heads or tails.

(7) The catch is (=) that the box has to follow a certain rule. Whatever you observe initially, (eb) you must observe the same thing ten seconds later.

(8) If you open the box and see heads, then (eb) ten seconds later when you open the box you must also see heads.

$$[a_7] = \frac{1K_7 b_7}{\delta_7 - b_7} + \frac{\sigma_7 \, \mathrm{Q}_7^*}{1 + \mathrm{Q}_7^*} > 1\vartriangle$$

This equation is important in its consequence in respect of following cases for which the model is applied:

**Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments**

(25) For the experiment we'll open the box, then open it five seconds later, then finally open it again at ten seconds. We know the first and last observations have to agree, but what about the observation in the middle? It depends upon what the box does with (e&eb) the coin.

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| (26) One possibility is that the box does nothing, so if you see tails the first time (eb), you will see tails both other times. |
| (27) Another possibility is that each time you open the box (eb) it flips the coin when you close it. |
| (28) So the first time you see heads, (eb) the second time tails, and (eb) the third time heads again. |
| (29) But they could also make a tricky box that flips (e&eb) the coin at seemingly random times, but does so in such a way that every ten seconds it repeats (eb) the pattern. In that case, authors would always see the same thing the first and last time, but the middle observation would seem (eb) random. |
| (30) If Macrorealism is true, then (eb) middle observations being always the same as (=) the first and last observation, always opposite or always random are (=) the only possibilities for the experiment |

**Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry Download as PDF (Size: 62KB) HTML PP. 50-53 DOI: 10.4236/jqis.2011.12007 5,508 Downloads 14,712 Views Citations Author(s) Leave a comment Mohamed S. El Naschie**

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| (25) Building upon the pioneering work of J. Bell [1] and an incredible result due to L. Hardy [2] it was shown (eb) that the probability of quantum entanglement of two particles is (=) a maximum of 9.0169945 percent [2]. |
| (26) This happens to be exactly the golden mean to (e) the power of five (?5) [3-7]. |
| (27) Although it has gone largely unnoticed for a long time, this result was essentially established independently in a much wider context by (e) the present author almost two decades ago [3-6]. |
| (28) The present work gives two fundamentally different derivations of (e) Hardy's beautiful result leading to (eb) precisely the same general conclusion, namely that by virtue of (e) the zero measure of the underlying |
| (29) Cantorian-fractal spacetime geometry the notion of (e) spatial separability in quantum physics is (=) devoid of any meaning [7]. The first derivation is purely logical and uses (e) a probability theory which combines the discrete with (e&eb) the continuum. |
| (30) The second derivation is purely geometrical and topological using (e) the fundamental equations of a theory developed by the author and his collaborators frequently referred to as E-infinity or Cantorian spacetime theory [3-7]. KEYWORDS Hardy's Quantum Entanglement, Golden mean, Cantor sets, Fractal Spacetime, E-Infinity Theory, Quantum Mechanics, J. S. Bell Cite this paper M. El Naschie, "Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry," Journal of Quantum Information Science, Vol. 1 No. 2, 2011, pp. 50-53. Doi: 10.4236/jqis.2011.12007. |

**Chaos, Solitons & Fractals Volume 28, Issue 4, May 2006, Pages 857–878 Sir Hermann Bondi 1919–2005 Cantorian spacetime and Hilbert space: Part I—Foundations G. Iovane http://dx.doi.org/10.1016/j.chaos.2005.08.074**

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| (41) Authors are going to show the link between the ε(∞) Cantorian space and( e&eb) the Hilbert spaces H(∞). |
| (42) In particular, El Naschie's ε (∞) is (=) a physical spacetime, i.e. an infinite dimensional fractal space, where (e) time is spacialized and the transfinite nature manifests (eb) itself. |
| (43) El Naschie's Cantorian spacetime is an arena where (e) the physics laws appear at each scale in a self-similar way linked to (e&eb) the resolution of the act of observation. |
| (44) By contrast the Hilbert space H (∞) is (=) a mathematical support, which describes (eb) the interaction between the observer and (e&eb) the dynamical system under (e&eb) measurement. |

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| (45) The present formulation, which is based on (e) the non-classical Cantorian geometry and topology of spacetime, automatically solves (e) the paradoxical outcome of (e) the two-slit experiment and the so-called particle–wave duality. |
| (46) In particular, measurement (i.e. the observation) is equivalent to (=) a projection of ε(∞) in the Hilbert space built on 3 + 1 Euclidean spacetime. |
| (47) Consequently, the wave–particle duality becomes (=) a mere natural consequence of (e) conducting an experiment in (eb) a spacetime with (e&eb) non-classical topological and (e&eb) geometrical structures, while (e) observing and taking measurements in (eb) a classical smooth 3 + 1 Euclidean spacetime. |
| (48) In other words, the experimental fact that a wave–particle duality exists is (=) an indirect confirmation of the existence of ε(∞) and a property of (e) the quantum–classical interface. |
| (49) Another direct consequence of the fact that real spacetime is) =) infinite dimensional hierarchical ε(∞) is (=) the existence of scaling law R(N), introduced by the author, which generalizes the Compton wavelength. |
| (50) It gives an answer to the problem of (e) segregation of matter at different scales, and shows (eb) the role of fundamental constants such as (=) the speed of light and Plank's constant h in the fundamental lengths scale without (e) invoking the principles of quantum mechanics. Copyright © 2005 Elsevier Ltd. All rights reserved |

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**Chaos, Solitons & Fractals Volume 14, Issue 9, December 2002, Pages 1341–1351 Fractal strings as an alternative justification for El Naschie's Cantorian spacetime and the fine structure constant Carlos Castro, http://dx.doi.org/10.1016/S0960-0779(02)00110-8**

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| (33) Beginning with the most general fractal strings/sprays construction recently expounded in the book by Lapidus and Frankenhuysen, it is shown how the complexified extension of El Naschie's cantorian-fractal spacetime model belongs to (e&eb) a very special class of families of fractal strings/sprays whose scaling ratios are given by (e) suitable pinary (pinary, p prime) powers of the Golden Mean. |
| (34) They then proceed to show why the logarithmic periodicity laws in (eb) Nature are direct physical consequences of (e) the complex dimensions associated with (e&eb) these fractal strings/sprays. |
| (35) Authors proceed with a discussion on quasi-crystals with (e&eb) p-adic internal symmetries, von Neumann's Continuous Geometry, the role of wild topology in (e&eb) fractal strings/sprays, (e&eb) the Banach-Tarski paradox, and tesselations of (e) the hyperbolic plane, quark confinement and (e&eb) the Mersenne-prime hierarchy of (e) bit-string physics in determining the fundamental physical constants in Nature. Fax: +1-914-965-3397 Copyright © 2002 Published by Elsevier Ltd. About ScienceDirectRemote access Shopping cart Contact and support Terms and conditions Privacy policy Cookies are used by this site. For more information, visit the cookies page. Copyright © 2016 Elsevier B.V. or its licensors or contributors. Science Direct ® is a registered trademark of Elsevier B.V. |

**Chaos, Solitons & Fractals Volume 21, Issue 2, July 2004, Pages 261–271 On Cantorian spacetime over number systems with division by zero Jakub Czajko, http://dx.doi.org/10.1016/j.chaos.2003.12.046**

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| (36) Division by zero is proposed in order to open (eb) new horizons for (e) abstract mathematics and physics. |
| (37) Although the formerly forbidden operation may be of (e) little practical value, the feasibility of unrestricted operations is needed for (e) meaningful deployment of linear vector spaces, whose duals are (=) inverse transposes of their primary vector spaces. |
| (38) Both quantum mechanics and (e&eb) pan-geometry of multispatial hyperspace operate on (e&eb) vectors that belong to (e&eb) mutually dual linear vector spaces. |
| (39) The division by zero also may remove (e) many years old paradox related to (e&eb) multiplication of zeros and makes it possible to introduce (eb) the concept of bigroup as (=) an additive and multiplicative group at the same time. |
| (40) This operational extension to (e&eb) algebraic groups makes thus (eb) an infinite-dimensional, fractal Cantorian spacetime (=) the preferred, invertible, operationally unrestricted, abstract mathematical infrastructure for physics. Tel.: +1-925-681-2058 Copyright © 2003 Elsevier Science Ltd. All rights reserved. |

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The nontrivial equilibrium is stable.

Some of the examples in which stability analysis is of seminal importance are: These are the system to which the

model becomes applicable

**Observation of a Dissipation-Induced Classical to Quantum Transition J. Raftery, D. Sadri, S. Schmidt, H. E. Türeci, and A. A. Houck Phys. Rev. X 4, 031043 – Published 8 September 2014**

(447)Here, authors report the experimental observation of (e) a dynamical quantum phase transition in (eb) a strongly interacting open photonic system.

(448) The system studied, comprising a Jaynes-Cummings dimer realized on (eb) a superconducting circuit platform, exhibits (eb) a dissipation-driven localization transition.

(449)Signatures of the transition in the homodyne signal and photon number reveal (eb) this transition to be from (e) a regime of classical oscillations into (e&eb) a macroscopically self-trapped state manifesting revivals, (=) a fundamentally quantum phenomenon.

(450) This experiment also demonstrates (eb) a small-scale realization of a new class of quantum simulator, whose well-controlled (e&eb) coherent and dissipative dynamics is suited to (e) the study of quantum many-body phenomena out of (e) equilibrium. Received 10 December 2013 DOI:http://dx.doi.org/10.1103/PhysRevX.4.031043 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**Viewpoint: Fermionic Vortices Find their Dual John McGreevy, Department of Physics, University of California at San Diego, La Jolla, CA 92093-0319, USA June 27, 2016• Physics 9, 72Theoretical work reveals a surprising relationship between the physics of fermionic vortices and quantum electrodynamics. Figure caption expand figure APS/Alan Stonebraker Figure 1: The action of discrete symmetries on the fermionic particle (top) and its dual fermionic vortex (bottom).**

(451)Dirac fermions in 2+1 dimensions (two spatial dimensions plus time) have been found in (eb) graphene and on (e&eb) the surfaces of topological insulators (TIs).

(452) A trio of papers now shows that these particles can also be found in (eb) a dual description of quantum electrodynamics (QED) [1–3].

(453)This deceptively simple statement joins (e&eb) together several disparate lines of progress in theoretical physics: interacting topological insulators, (e&eb) anomalies in quantum field theory, (e&eb) spin liquids, and (e&eb) compressible Quantum Hall States.

(454)Humans are not very good at understanding interacting many-body systems.

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| (455)We understand the behavior of free particles and can sometimes use (e) perturbation theory to (e) incorporate weak interactions between (e&eb) them. |
| (456)But what happens when we couple (e&eb) the particles strongly? |
| (457)To reveal the interesting collective phenomena that can result, we've had to rely largely on (e) experiments. |
| (458)Examples include (e) superconductivity (e&eb) superfluidity (e&eb) the integer and (e&eb) fractional quantum Hall effects (e&eb) the fractionalization of charge and spin, and (e&eb) the wetness of water. |
| (459)The fractional quantum Hall effect (FQHE) is a physical phenomenon in which the Hall conductance of 2D electrons shows (eb) precisely quantised plateaus at (eb) fractional values of (e) {\displaystyle e^{2}/h} e^{2}/h. |
| (460) It is a property of (e) a collective state in which electrons bind (e&eb) magnetic flux lines to make (eb) new quasiparticles, and excitations have (e) a fractional elementary charge and possibly also fractional statistics. The 1998 Nobel Prize in Physics was awarded to Robert Laughlin, Horst Störmer, and Daniel Tsui for the discovery and explanation of the fractional Hall effect.[1] |
| (461) However, Laughlin's explanation was (=) a phenomenological guess and only applies to (e&eb) fillings {\displaystyle \nu =1/m} \nu =1/m where (e) {\displaystyle m} m is an odd integer. The microscopic origin of the FQHE remains unknown and is a major research topic in condensed matter physics. (Wikipedia) |
| (462)The quantum Hall effect (or integer quantum Hall effect) is a quantum-mechanical version of (e) the Hall effect, observed in (eb) two-dimensional electron systems subjected to (e&eb) low temperatures and strong magnetic fields, in which the Hall conductance σ undergoes (e&eb) certain quantum Hall transitions to take on (e&eb) the quantized values {\displaystyle \sigma ={\frac {I_{\text{channel}}}{V_{\text{Hall}}}}=\nu \;{\frac {e^{2}}{h}},} \sigma = \frac{I_\text{channel}}{V_\text{Hall}} = \nu \; \frac{e^2}{h}, where {\displaystyle I_{\text{channel}}} I_\text{channel} is the channel current, {\displaystyle V_{\text{Hall}}} V_\text{Hall} is the Hall voltage, e is the elementary charge and h is Planck's constant. The prefactor ν is known as the "filling factor", and can take on either integer (ν = 1, 2, 3, ...) or fractional (ν = 1/3, 2/5, 3/7, 2/3, 3/5, 1/5, 2/9, 3/13, 5/2, 12/5, ...) values. |
| (463) The quantum Hall Effect is referred to as (=) the integer or fractional quantum Hall effect depending on whether ν is an integer or fraction, respectively. |
| (464)The striking feature of the integer quantum Hall Effect is (=) the persistence of the quantization (i.e. the Hall plateau) as (e) the electron density is varied. |
| (465) Since the latter remains constant when (e) the Fermi level is in a clean spectral gap, it must be that the Fermi level is immersed in (eb) essential spectrum (i.e. of finite density of states), though Anderson localized. |
| (466) The quantization and stability of the Hall conductance in such conditions was achieved in 1994 by (e) Bellissard, van Elst and Schulz-Baldes using the methods of (e) Non-Commutative Geometry.[1] |
| (467)The fractional quantum Hall effect is more complicated, as its existence relies fundamentally on (e) electron–electron interactions. |
| (468)Although the microscopic origins of the fractional quantum Hall effect are unknown, there are several phenomenological approaches that provide (eb) accurate approximations. |
| (469) For example the effect can be thought of as (=) an integer quantum Hall effect, not of electrons but of (e) charge-flux composites known as (=) composite fermions. |

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| (470) In 1988, it was proposed that there was quantum Hall effect without Landau levels. This quantum Hall Effect is referred to as (=) the quantum anomalous Hall (QAH) effect. |
| (471)There is also a new concept of the quantum spin Hall effect which is an analogue of the quantum Hall effect, where (e) spin currents flow instead of (e&eb) charge currents.[2] |
| (472)In two dimensions, when classical electrons are subjected to (e&eb) a magnetic field they follow (eb) circular cyclotron orbits. |
| (473)When the system is treated quantum mechanically (eb) these orbits are quantized. |
| (474)The energy levels of these quantized orbitals take on (e) discrete values |
| (475):{\displaystyle \textstyle E_{n}=\hbar \omega _{c}(n+1/2),} \textstyle E_{n}=\hbar \omega _{c}(n+1/2), where ωc = eB/m is the cyclotron frequency |
| (476) These orbitals are known as Landau levels, and at weak magnetic fields, their existence gives rise to (eb) many interesting "quantum oscillations" such as the Shubnikov–de Haas oscillations and the de Haas–van Alphen effect (which is often used to map the Fermi surface of metals). |
| (477) For strong magnetic fields, each Landau level is (=) highly degenerate (i.e. there are many single particle states which have the same energy En). |
| (478)Specifically, for a sample of area A, in magnetic field B, the degeneracy of each Landau level is (=) {\displaystyle \textstyle N=g_{s} BA/\phi _ {0},} \textstyle N=g_{s} BA/\phi _ {0}, where gs represents (eb) a factor of 2 for spin degeneracy, and ϕ0 = 2·10−15 Wb is the magnetic flux quantum. |
| (479) For sufficiently strong B-fields, each Landau level may have (e) so many states that all of the free electrons in the system sit in only a few Landau levels; it is in this regime where (e) one observes the quantum Hall Effect. |
| (480)The integers that appear in the Hall Effect are (=) examples of **topological quantum numbers**. |
| (481) They are known in mathematics as the first **Chern numbers** and are closely related to (e&eb) **Berry's phase.** |
| (482)A striking model of much interest in this context is the Azbel-Harper-Hofstadter model whose quantum phase diagram is (=) the Hofstadter butterfly shown in the figure. (deleted due to spatial constraints) |
| (483)The vertical axis is (=) the strength of the magnetic field and the horizontal axis is the chemical potential, which fixes the electron density. The colors represent the integer Hall conductance. Warm colors represent positive integers and cold colors negative integers. The phase diagram is fractal and has structure on all scales. In the figure there is an obvious self-similarity. |
| (484)Concerning physical mechanisms, impurities and/or particular states (e.g., edge currents) are important for (e) both the 'integer' and 'fractional' effects. |
| (485) In addition, Coulomb interaction is also essential in (eb) the fractional quantum Hall effect. |
| (486)The observed strong similarity between integer and (=) fractional quantum Hall effects is explained by (e) the tendency of electrons to form bound states with (e&eb) an even number of magnetic flux quanta, called composite fermions. |
| (487)Duality is the circumstance in which cranking up interactions results in (eb) another system we understand, usually one that can be described in terms of (e&eb) another set of weakly coupled particles. |
| (488)When such a dual representation is found, (eb) it gives a solution to the strong coupling problem. |
| (489) A further reason that theoretical physicists get excited about dualities is that they undermine (e) the very notion of an “elementary particle.” |
| (490)In terms of the original variables, the dual, weakly coupled particles are (=) generally solitons—large, heavy, classical objects made up of (e) many quanta. |

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| (491)The remarkable thing is that solitons can become (=) "elementary," in the sense that the model that describes the system can be reformulated with the solitons as (=) the basic constituents. |
| (492)An example of a duality, which we will use below, arises in (eb) Maxwell's theory of electromagnetism, where (e) the strength of the interactions is determined by (e) the electric charge e. |
| (493)The duality exchanges electric and (e&eb) magnetic point charges, and the interaction strength in the dual description is 1⁄e, which is small when (e) e is large. |
| (494)This idea has been generalized to (e) many other field theories, thanks, in particular, to tools of (e) supersymmetry [4]. |
| (495)Another classic example of duality involves (e&eb) vortices. |
| (496) A vortex is a localized object in the plane around which a phase field winds (e&eb). (A phase field takes values on the circle.) |
| (497)Dualities relating particles and vortices have (e) a rich history in particle physics and condensed-matter physics. |
| (498) The simplest version relates the particles in the normal state of a superfluid to (e&eb) the vortices in a dual superconductor [5]. |
| (499)This charge-vortex duality has been (=) a workhorse in the study of two-dimensional condensed-matter systems, such as (=) fractional Quantum Hall States, superconducting thin films, and beyond-Landau critical phenomena [6]. |
| (500)Vortices can be (=) fermions (as they frequently are, for instance, in type-II superconductors). |
| (501) A dual description of a system with (e&eb) fermionic vortices thus requires (e) a new ingredient, since (e) it must include (e) such a fermionic particle. |
| (502)The new work [1–3] builds on (e&eb) earlier theoretical work [7] to make (eb) an appealing proposal for (e) the required duality: the fermionic vortex of QED in (eb) 2+1 dimensions is (=) a free Dirac fermion. The argument for this proposal builds on our understanding of TIs. |
| (503)TIs are distinguished only by (e) their edge states. |
| (504) The most famous example is a (3+1)-dimensional time-reversal invariant TI, which can host boundary states whose band structure is defined by (e) a single Dirac cone. |
| (505)A single time-reversal-symmetric Dirac cone is not possible without (e) the bulk TI, so its appearance in experiments provides (eb) a signature of the bulk phase. |
| (506) But the Dirac cone is not (e) the only possible signature. |
| (507) Deformations of the boundary conditions can dramatically change (e&eb) the spectrum of the edge states, yet preserve some essential "TI-ness." This essential property is (=) called an anomaly in high-energy theory. |
| (508)Max Metlitski of the University of California, Santa Barbara, and Ashvin Vishwanath of the University of California, Berkeley [1], and independently, Chong Wang and Todadri Senthil at the Massachusetts Institute of Technology, Cambridge [2], use this ambiguity in the edge theory and the electric-magnetic duality of an auxiliary theory of bulk electromagnetism coupled to (e&eb) the TI to argue for the duality. |
| (509)When I said that (3+1)-dimensional QED had a duality that interchanged the electric and magnetic charges, one could have (e) complained that the electron is (=) a fermion, while (e) a magnetic monopole may not be. |
| (510)However, the authors observe that in a TI, a charge-two monopole is (=) also a fermion. |

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| (511) By analyzing the spectrum of electric and magnetic charges, they conclude that the duality exchanges the electron and (e&eb) the charge-two monopole and, moreover, exchanges a TI with (e&eb) time-reversal symmetry with (e&eb) a "chiral TI," which instead has (e) particle-hole symmetry |
| (512). But this is just what is required of (e) a symmetry that exchanges a fermionic particle with (e&eb) a fermionic vortex: time-reversal symmetry maps the electron to itself but maps the vortex to (e&eb) another vortex whose phase winds (e&eb) the other way (see Fig. 1). |
| (513)The exchange of particles and holes, in turn, preserves (eb) the vortex. |
| (514)The bulk duality implies (eb) a duality between edge theories. |
| (515)The freedom to deform the edge of the TI implies (eb) a weak form of equivalence: although the edge theories are (=) topologically the same (in terms of symmetry assignments and (e&eb) anomalies), they need not have (e) identical low-energy physics. |
| (516)One side of the duality is (=) a free, massless Dirac particle, but the long-distance physics of (2+1)-dimensional QED is (=) a long-standing problem, which is important in the study of spin liquids. |
| (517) Many gauge theories in two dimensions develop (eb) a gap in their energy spectrum because of (e) the proliferation of monopoles. |
| (518) In the realization proposed by Metlitski and Vishwanath [1] and by Wang and Senthil [2], this cannot happen because (e) the monopoles carry electric charge, and indeed the stronger notion of duality as (=) an equivalence of low-energy physics may hold. |
| (519)Evidence for the stronger relation comes from (e) David Mross and colleagues at the California Institute of Technology, Pasadena [3], who have derived a lattice version of the duality using (e) a clever but explicit change of variables. |
| (520)Their construction is particularly interesting because (e) it uses (e) a lattice of one-dimensional models which themselves can only occur at the edge of a TI. |
| (521)Authors explicitly derive the duality between a free electronic Dirac cone and (e&eb) quantum electrodynamics in (eb) (2+1) dimensions (QED3) with (e&eb) N=1 fermion flavors. |
| (522)The duality proceeds via (e&eb) an exact, nonlocal mapping from electrons to (e&eb) dual fermions with (e&eb) long-range interactions encoded by (e) an emergent gauge field. |
| (523)This mapping allows (eb) us to construct (eb) parent Hamiltonians for exotic topological-insulator surface phases, derive (eb) the particle-hole-symmetric field theory of (e) a half-filled Landau level, and nontrivially constrain (e) QED3 scaling dimensions. |
| (524)They similarly establish duality between bosonic topological insulator surfaces and (e&eb) N=2 QED3. |

**Black hole thermodynamics with generalized uncertainty principle Li Xiang, X. Q. Wen (Submitted on 6 Jan 2009 (v1), last revised 18 Mar 2009 (this version, v2))**

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| (525) In the standard viewpoint, the temperature of a stationary black hole is proportional to (=) its surface gravity, $T_H=\hbar\kappa/2\pi$. |
| (526)This is a semiclassical result and the quantum gravity effects are not (e) taken into consideration. |
| (527)This Letter explores a unified expression for (e) the black hole temperature in the sense of (e) a generalized uncertainty principle (GUP). |
| (528) Discussion involves a heuristic analysis of (e) a particle which is absorbed by (e) the black hole. |
| (529)Besides a class of static and spherically symmetric black holes, an axially symmetric Kerr-Newman black hole is considered. Different from the existing literature, authors suggest (eb) that the black hole's irreducible mass represent (eb) the characteristic size in (eb) the absorption process. |

(530)The information capacity of a remnant is also discussed by (e) Bousso's D-bound in de Sitter spacetime. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP 0910:046,2009 DOI: 10.1088/1126-6708/2009/10/046 Cite as: arXiv:0901.0603 [gr-qc] (or arXiv:0901.0603v2 [gr-qc] for this version)

IF we define

$$[a_7]_7 = [f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}]\,P_7^*$$

$$[a_7]_7 = \left[f_{7*}^{7}\right] P_7^* - \frac{h_*^7}{h_*} P^*(0)$$

$$e = [g_*^7]_7\, Q_7^*$$

$$\eta_7(t) = \left[h_*\right]_7 \beta_7 \left(\theta_7(t)\right) - \mu_7(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\left[a_7\right] M_7 - [a_7]_7 W_7 - [n(t)] M_7$$

$$w_7(t) = eM_7$$

Write

$$A = \begin{pmatrix} -[a_7]_7 - [a_7]_7 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B = (t) = \begin{pmatrix} [\eta(t)]_7 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_7 \\ W_7 \end{pmatrix}$$

The solution of (32)is

$$X_7(t) = e^{A,t}[X]_7(0) + \int_0 e^{At}$$

## PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION

## ATTRITION MODELS

## MODULE EIGHT

Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal.

### 43. MODEL FORMULATION

**Model is applied to systems given in the following:**

**Researchers blur the line between classical and quantum physics by connecting chaos and entanglement July 12, 2016 by Sonia Fernandez**

(2599) Experimental link between quantum entanglement (left) and (e&eb) classical chaos (right) found using (e) a small quantum computer. Credit: University of California - Santa Barbara Using a small

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| quantum system consisting of three superconducting qubits, researchers at UC Santa Barbara and Google have uncovered a link between aspects of classical and (e&eb) quantum physics thought to be unrelated: classical chaos and (e&eb) quantum entanglement. |
| (2600) Their findings suggest (eb) that it would be possible to use (e) controllable quantum systems to investigate certain fundamental aspects of nature. |
| (2601) "It's kind of surprising because (e) chaos is this totally classical concept—there's (=) no idea of chaos in a quantum system," Charles Neill, a researcher in the UCSB Department of Physics and lead author of a paper that appears in Nature Physics. |
| (2602) "Similarly, there's no concept of entanglement within (eb) classical systems. And yet it turns out that chaos and (e&eb) entanglement are really very strongly and clearly related." |
| (2603) Initiated in the 15th century, classical physics generally examines and describes (eb) systems larger than atoms and molecules. |
| (2604) It consists of hundreds of years' worth of study including Newton's laws of motion, electrodynamics, relativity, thermodynamics as well as chaos theory—the field that studies the behavior of (e) highly sensitive and unpredictable systems. |
| (2605) One classic example of chaos theory is the weather, in which a relatively small change in one part of (e) the system is enough to foil predictions—and vacation plans—anywhere on the globe. |
| (2606) At smaller size and length scales in nature, however, such as those involving (e&eb) atoms and photons and their behaviors, classical physics falls short. |
| (2607) In the early 20th century quantum physics emerged, with its seemingly counterintuitive and sometimes controversial science, including (e) the notions of superposition (the theory that a particle can be located in (eb) several places at once) and entanglement (particles that are deeply linked behave as (=) such despite physical distance from one another). |
| (2608) And so began the continuing search for connections between the two fields. All systems are fundamentally quantum systems, according Neill, but the means of describing in a quantum sense the chaotic behavior of, (e) say, air molecules in an evacuated room, remains limited. |
| (2609) Imagine taking a balloon full of air molecules, somehow tagging them so you could see them and then releasing them (eb) into a room with no air molecules, noted co-author and UCSB/Google researcher Pedram Roushan. |
| (2610) One possible outcome is that the air molecules remain clumped together (e&eb) in a little cloud following the same trajectory around the room |
| (2611) And yet, he continued, as we can probably intuit, the molecules will more likely take off in (eb) a variety of velocities and directions, bouncing off walls and interacting with (e&eb) each other, resting after the room is (=) sufficiently saturated with them. |
| (2612) "The underlying physics is (=) chaos, essentially," he said. |
| (2613) The molecules coming to rest—at least on (eb) the macroscopic level—is the result of (e) thermalization, or of reaching equilibrium after they have achieved (eb) uniform saturation within the system. |
| (2614) But in the infinitesimal world of quantum physics, there is (=) still little to describe that behavior. |
| (2615) The mathematics of quantum mechanics, Roushan said, do not allow for (e) the chaos described by Newtonian laws of motion. |
| (2616) To investigate, the researchers devised an experiment using (e) three quantum bits, the basic computational units of (e) the quantum computer. |

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| (2617) Unlike classical computer bits, which utilize (e) a binary system of two possible states (e.g., zero/one), a qubit can also use (e) a superposition of both states (zero and one) as a single state. |
| (2618) Additionally, multiple qubits can entangle, or link so closely that their measurements will (=) automatically correlate. |
| (2619) By manipulating these qubits with electronic pulses, Neill caused them to interact, rotate and evolve (e&eb, e&eb, eb) in the quantum analog of (e) a highly sensitive classical system. |
| (2620) The result is a map of entanglement entropy of (e) a qubit that, over time, comes to strongly resemble that of (e) classical dynamics—the regions of entanglement in (eb) the quantum map resemble the regions of (e) chaos on the classical map. |
| (2621) The islands of low entanglement in (eb) the quantum map are located in the places of (e) low chaos on the classical map. |
| (2622) "There's a very clear connection between entanglement and (e&eb) chaos in these two pictures," said Neill. |
| (2623) "And, it turns out that thermalization is (=) the thing that connects chaos and (e&eb) entanglement. |
| (2624) It turns out that they are actually the driving forces behind (eb) thermalization. |
| (2625) "What we realize is that in almost any quantum system, including on quantum computers, if you just let it evolve and you start to study what happens as a function of time, (eb) it's going to thermalize," added Neill, referring to the quantum-level equilibration. "And this really ties together the intuition between classical thermalization and (e&eb) chaos and how it occurs (eb) in quantum systems that entangle." |
| (2626) The study's findings have fundamental implications for (e) quantum computing. |
| (2627) At the level of three qubits, the computation is relatively simple, said Roushan, but as researchers push to build increasingly sophisticated and powerful quantum computers that incorporate (e) more qubits to study highly complex problems that are beyond the ability of classical computing—such as those in the realms of machine learning, artificial intelligence, fluid dynamics or chemistry—a quantum processor optimized for such calculations will be a very powerful tool. |
| (2628) "It means we can study things that are completely impossible to study right now, once we get to bigger systems," said Neill.More information: Nature Physics, DOI: 10.1038/nphys3830 , http://www.nature.com/nphys/journal/vaop/ncurrent/full/nphys3830.html Journal reference: Nature Physics search and more info website Provided by: University of California - Santa Barbara |

**Title: coherent control of macroscopic quantum states in a single-Cooper-pair box Authors: Nakamura, Y.; Pashkin, Yu. A.; Tsai, J. S. Publication: Nature, Volume 398, Issue 6730, pp. 786-788 (1999). (Nature Homepage) Publication Date: 04/1999 Origin: NATURE DOI: 10.1038/19718 Bibliographic Code: 1999Natur.398...786N**

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| (2629) A nanometre-scale superconducting electrode connected to (e&eb) a reservoir via (e&eb) a Josephson junction constitutes (e) an artificial two-level electronic system: a single-Cooper-pair box. |
| (2630) The two levels consist of (e) charge states (differing by 2e, where (e) e is the electronic charge) that are coupled by (e) tunnelling of Cooper pairs through (e&eb) the junction. |
| (2631) Although the two-level system is macroscopic, containing (e) a large number of electrons, the two charge states can be (=) coherently superposed. |

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| (2632) | The Cooper-pair box has therefore been suggested as (=) a candidate for a quantum bit or `qubit'- the basic component of a quantum computer. |
| (2633) | Here authors report the observation of quantum oscillations in (eb) a single-Cooper-pair box. |
| (2634) | By applying a short voltage pulse via (e&eb) a gate electrode, we can control (e&eb) the coherent quantum state evolution: the pulse modifies (e&eb) the energies of (e) the two charge states non-adiabatically, bringing them into (e&eb) resonance. |
| (2635) | The resulting state-a superposition of (e) the two charge states-is detected by (e) a tunnelling current through (e&eb) a probe junction. |
| (2636) | Results demonstrate electrical coherent control of (e) a qubit in a solid-state electronic device. |

**Superconducting quantum computing From Wikipedia, the free encyclopedia**

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| (2637) | Superconducting quantum computing is a promising implementation of (e) quantum information technology that involves (e&eb) nanofabricated superconducting electrodes coupled through (e&eb) Josephson junctions. |
| (2638) | As in a superconducting electrode, the phase and (e&eb) the charge are conjugate variables. |
| (2639) | There exist three families of superconducting qubits, depending on (=) whether the charge, (e&eb) the phase, or neither of (e) the two is good quantum number. |
| (2640) | These are respectively termed charge qubits (e&eb) flux qubits, and (e&eb) hybrid qubits. |

**Theory**

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| (2641) | Unlike many other physical implementations of a qubit which involve (e&eb) exclusively two-level systems (such as nuclear spin and (e&eb) photon polarization), the integrated quantum circuit involved in (e&eb) a superconducting qubit is (=) a multi-level system, of which only the first two levels are used as (=) the computational basis. |
| (2642) | A basic requirement for such an implementation is that the energy levels are (=) not uniformly spaced, so that (e) photons of a particular frequency which cause (eb) transition between the 0 and 1 levels do not(e) cause transitions from the first level to (e&eb) the higher levels as well. |
| (2643) | For electronic signals to be carried from one part of the circuit to (e&eb) another without energy loss and (e&eb) hence decoherence, the system also needs to be (=) non-dissipative, i.e., the metallic parts involved should have (e) zero resistance. |
| (2644) | These circuits currently need to be (=) operated at very low temperatures so that, 1) superconductivity is (=) realized, and 2) thermal fluctuations do not (e) cause transitions between (e&eb) energy levels. |

**Quantum LC Circuit**

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| (2645) | The simplest non-dissipative quantum circuit consists simply of (e) an inductor L {\displaystyle L} L and (e&eb) capacitor C {\displaystyle C} C, with (e&eb) the metallic wires connecting them being superconducting. |
| (2646) | Here the flux Φ {\displaystyle \Phi } \Phi through (e&eb) the inductor and (e&eb) the charge Q {\displaystyle Q} Q in the capacitor are (=) canonically conjugate variables which obey (eb) the commutation relation [ Φ , Q ] = i ℏ {\displaystyle [\Phi ,Q]=i\hbar } [\Phi ,Q]=i\hbar |

**The Hamiltonian associated with this system is**

(2647) H = Φ 2 2 L + Q 2 2 C {\displaystyle H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}} H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}

(2648) This gives rise to the energy levels E = ℏ ω 0 ( n + 1 2 ) {\displaystyle E=\hbar \omega _{0}(n+{\frac {1}{2}})} E=\hbar \omega _{0}(n+{\frac {1}{2}}) Here ω 0 = 1 L C {\displaystyle \omega _{0}={\frac {1}{\sqrt {LC}}}} \omega _{0}={\frac {1}{{\sqrt {LC}}}}

**Quantum gravity Carlo Rovelli (2008), Scholarpedia, 3(5):7117. doi:10.4249/scholarpedia.7117 revision #135539 [link to/cite this article] Post-publication activity Curator: Carlo Rovelli Dr. Carlo Rovelli, Center of Theoretical Physics, and Marseille, France**

(2649) Quantum Gravity is the name given to any theory that describes gravity in (eb) the regimes where quantum effects cannot be disregarded.

(2650) At present, there is no such a theory which is universally accepted and confirmed by experience. Therefore the term "Quantum Gravity" indicates (eb) more an open problem than a specific theory.

(2651) Several research lines, at different levels of development, offer tentative solutions to the problem. These tentative quantum-gravity theories are variously viewed as competing research directions, or as contributions to the common goal of finding the physically correct theory. The quest for the good quantum theory of gravity bears on a number of fundamental issues and it is sometimes presented as the most important open problem in fundamental physics: the "Holy Grail" of contemporary theoretical physics.

**The problem of quantum gravity**

(2652) The current theory of (e) gravity is general relativity, a theory developed by Albert Einstein at the beginning of the XX century, which has changed in depth our understanding of (e) what are space and time.

(2653) Similarly, quantum mechanics has changed in depth our understanding of matter, energy and causality. These two theories, today very well empirically supported, have opened (eb) a major conceptual revolution in physics.

(2654) However, they appear to be incompatible, at least at a first reading, because (e) each of the two is formulated on the basis of (e) principles that are explicitly contradicted by (e) the other theory.

(2655) Therefore in spite of the immense development of scientific knowledge it has delivered, the physics of the XX century has left us with (e&eb) a major confusion on what are the basic conceptual ingredients for (e) understanding the physical world.

(2656) In other words, the scientific revolution opened by general relativity and quantum mechanics at the beginning of (e) the XX century is not concluded yet, and a new synthesis is required. Quantum gravity, merging general relativity and quantum mechanics should be this new synthesis.

**Quantum space and quantum time**

(2657) Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.

(2658) In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.

(2659) They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.

(2660) When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.

(2661) Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.

**The Planck scale**

(2662) Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".

(2663) Here $\hbar$ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.

(2664) The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.

(2665) Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)

(2666) Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.

## Physical applications

Some phenomena where quantum gravity plays a major role are the following.

**The microscopic structure of spacetime**

(2667) If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.

(2668) Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length

(2669) Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.

(2670) A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.

**Early cosmology**

(2671) According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.

(2672) However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.

(2673) A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.

(2674) Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".

(2675) One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.

**Black holes**

(2676) Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.

(2677) First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.

(2678) Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.

(2679) What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.

(2680) Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.

**Astrophysical effects**

(2681) Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.

(2682) An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.

(2683) The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.

**Main approaches**

(2684) The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.

(2685) Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.

(2686) The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.

**String theory**

(2687) String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.

(2688) Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.

(2689) String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.

(2690) Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.

(2691) However, when summed the entire series appears to be (=) divergent.

(2692) In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.

(2693) Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.

(2694) Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.

(2695) In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.

**Loop quantum gravity and spin foams**

(2696) Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.

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| (2697) The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity. |
| (2698) Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime. |
| (2699) The granular properties of space can be (=) explicitly computed. |
| (2700) In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory. |
| (2701) The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks. |
| (2702) Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure. |
| (2703) The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit. |
| (2704) Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism. |
| (2705) This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism. |

**Noncommutative geometry**

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| (2706) Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime. |
| (2707) Quantum mechanics teaches us that dynamical quantities are (=) noncommutative --meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed. |
| (2708) It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry. |
| (2709) A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces. |
| (2710) Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces. |
| (2711) This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity. |

**Others**

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| (2712) Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays |

(2713) . Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection.

(2714) Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible.

(2715) The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.

Let $p_8(t, a_8)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(g_8)Q_8$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_8(0) = 0, \quad g_8(P_8) > 0, \quad P_8 \ 1, \qquad \lim_{P_8 \to \infty} g_8(P_8) = (K_8)_8 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_8 + (\rho_a)_8 + D_8(a_8, P_8, Q_8)\rho_8 + 0, \qquad 0 < a_8 \ \infty,$$

$$\rho_8(t,0) = \int_0^\infty n_8(a_8, P_8, Q_8)\rho_8(t_8, a_8)da_8$$

Where $D_8$ is the mortality function and $n_8$ is the fertility function.

Put $_8$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_8}{dt}\right]_8 = x_8\left[\frac{dQ_8}{dt}\right]_8 - x_8\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_8}{dt}\right]_8 = x_3\left[\frac{dQ_3}{dt}\right]_8 - x_4\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_8}{dt}\right]_8 - x_6\left[\frac{dQ_3}{dt}\right]_8$$

Where $x_8, x_8, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_8$ and $g_8(P_8)Q_8$

Now when you add the term $(b_8Q_8)$ with $g_8(P_8)(Q_8)$ we are infect see is that $(b_8Q_8)$ is credited incrementally part of $g_8(P_8)(Q_8)$ or instant nearly, when subtracted from $\frac{dQ_8}{dt}$ , it means that $g_8(P_8)(Q_8)$ is slowly eaten by $\frac{dQ_8}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_8 = \mu_8(a_8) + (f_8)_8(P_8)_8 + (f_8)_8(Q)_8 \quad (2)$$

$$n_8 = \beta_8(a_8)h_8(Q_8), \quad (3)$$

Where $\mu_8$the age- dependent debit rate whit is $\mu_8(a_8)$ 0 $\forall a > 0$; $(f_8)_8$ is the density- dependent death rate with $f_8(0) = 0$, $(f_8^8)_8(P_8)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_8)_8(P_8) = (k_8)_8$ ; $f_8$ is the functional response of creditors to prey with $[f_8]_8(0) = 0, (f_8^8)_8(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_8)_8(Q_8) = [k_8]$ ; $\beta$ is the age

dependent birth rate with $\beta(a) \geq 0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^8(Q) \geq 0, \forall Q \geq 0$, and $\lim_{Q\to\infty} h(Q) = k_3 \geq 0$.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_8 + (\rho_a)_8 = -\mu_8(a_8) + (f_8)_8 P_8) + \big(f_8)_8(Q_8)\big)(\rho_8)_8 \quad (4)$$

$$\rho_8(t_8, 0) = h_8(Q_8) \int_0^\infty \beta_8(a_8)\rho_8(t, a_8)da, \quad (5)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(P_8)Q_8 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_8$ and $Q_8$ of (4-6) satisfy

$$0 \leq \rho_8(t, a_8) \leq \rho_8^*(t, a_8),$$

$$0 \leq Q_8(t) \leq Q_8^*(t),$$

Where $\rho_8^*(t, a_8)$ and $Q_8^*(t)$, are solutions of

$$(\rho_t^*)_8 + (\rho_a^*)_8 = -(\mu_8(a_8) + (f_8)_8(P_8^*) + (f_8)_8 Q_8^*))\,\rho_8^*, \quad (7)$$

$$\rho_8^*(t, 0) = \int_0^\infty \beta_8(a_8)\,\rho_8^*(t, a_8)da_8 \quad (8)$$

$$\frac{dQ_8^*}{dt} = -b_8 Q_8^* + g_8(P_8^*)Q_8^*, \quad (9)$$

With $P_8^*(t) = \int_0^\infty \rho_8^*\ (t, a_8)da_8$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_8^*(t)}{dt} = P_8^*(t)\big(r_8 - (f_8)(P_8^*) - (f_8)_8(Q^*)\big), \tag{10}$$

$$\frac{dQ_8^*(t)}{dt} = Q_8^*(t)(-b_8 + g_8(P_8^*)), \tag{11}$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_8(a)e^{-\int_0^a \mu_8(\tau)dt}e^{-ra}\ da_8 = 1. \tag{12}$$

If $r < 0$, the trivial solution $P_8^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(t)dt}\ da_8 > 1, \tag{$H1$}$$

Which guarantees $r_8 > 0$.

*Remark.* If $\beta_8(a) = [\beta_0]_8 a_8 e^{-\alpha_8 a_8}, \mu_8(a_8) \equiv \mu_8$ then the solution of (12) is

$$r_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_8]_8(P_8) = 0, [f_8](Q_8) = \lambda_8 Q_8\ h_8(Q_8) = 1, \qquad and\ \ g_8(P_8) = c_8 P_8$$

$$[\rho_8]_8 + [\rho_a]_8 = -(\mu_8 + \lambda_8 Q_8)\rho_8$$

$$\rho_8(t,0) = \int_0^{\infty} [\beta_0]_8 a_8 e^{-\alpha_8 a_8} \rho_8(t,a_8)\, da_8$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + c_8 P_8 Q_8$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_8}{dt} = \omega_8 P_8 - \lambda_8 P_8 Q_8$$

$$\frac{dQ_8}{dt} = b_8 Q_8 - c_8 P_8 Q_8$$

Where $\omega_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$

## 44. THE EQUILIBRIUM $(\rho_0)_8(a_8), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_8(a_8), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_8 + (\rho_a)_8 = -\big(\mu_8(a_8) + (f_8)_8(P_8)\big)\rho,$$

$$\rho_8(t,0) = \int_0^{\infty} \beta_8(a_8)\rho_8(t,a_8) da_8 \qquad (13)$$

The solution of (13) has the form of

$$\rho_8(t,a_8) = Ce_8^{r(t-a_8)} - \int_0^a (\tau)d\tau - \int_0^8 (f_8)_8\big(P_8(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_8(t) = C_8 e^{r_8 t - \int_0^t (f_8)_8(P_8(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_8^{-8}(r)$ if $k_8 > r$. This is equivalent to

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{(k_8)_8 a_8} da_8 < 1.$$

In fact, if we define

$$R_8(s_8) = \int_0^\infty \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau)d\tau} e^{-s_8 a_8}\, da_8$$

$R(s)$ is a decreasing function. $R_8(r_8) = 1$ And $R_8(k_8)_8 < 1$ imply $r_8 < (k_8)_8$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_8(a_8), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{k_8 a_8} da_8 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_8 \equiv \int_0^\infty (\rho_0)_{88}(a_8)da_8$ *satisfies*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{-(f_8)_8(P_0)a_8} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_8(a_8), 0)$

Let $u_8(t,a_8) = \rho(t,a_8) - \rho_0(a_8)$ Then

$$N_8(t) \equiv \int_0^\infty u_8(t,a_8)da_8 = \int_0^\infty \rho_8(t,a_8)da_8 - \int_0^\infty (\rho_0)_8(a_8)da = P_8(t) - P_0)_8$$

And

$$(u_8)_8 + (u_a)_8 = (\rho_t)_8 + (\rho_a)_8 - (\rho_0^8)_8(a_8)$$

$$= -\mu_8(a_8) + (f_8)_8(P_8) + (f_8)_8(Q_8))\,(u_8 + (\rho_0)_8) + \mu_8(a_8)(\rho_0)_8 + (f_8)_8(P_0)_8(\rho_0)_8$$

$$= -\mu_8(a_8)u_8 - (f_8)_8(P_8) + (f_8)_8(Q_8))\,u_8 - (\rho_0)_8\,(f_8)_8(P_8) - (f_8)_8(P_0)_8$$

$$+ (f_8)_8(Q_8)), \qquad (15)$$

Where $(\rho_0^8)_8(a_8)$ denotes $d(\rho_0)_8(a_8)/da_8$

Expansion $(f_8)_8(P_8)$ and $(f_8)_8(Q_8)$ at $(P_0)_8$ and 0 yield respectively

$$(f_8)_8(P_8) = (f_8)_8((P_0)_8 + N_8)$$

$$= (f_8)_8(P_0)_8 + (f_8^8)_8(P_0)_8N_8 + o_8(N_8), \qquad (16)$$

$$(f_8)_8(Q_8)$$

$$= (f_8^8)_8(0)Q + o_8(Q_8) \qquad (17)$$

Write $f_8(P_0) = f_{80}, f_8'(P_0) = f_{80}'$, and $f_8'(0) = f_{80}'$ and substitute (16) and (17) into (15) to obtain

$$(u_8)_8 + (u_a)_8 \approx -(\mu_8(a_8) + (f_{80})_8)u - (f_{80}^8)_8(\rho_0)_8(N_8) - (f_{80}^8)_8(\rho_0)_8(Q_8)$$

$$u_8(t,0) = \rho_8(t,0)(\rho_0)_8(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a) + (\rho_0)_8(a_8)da_8 - (\rho_0)_8\,(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a_8)da_8 + (h_8(Q_8) - 1)(\rho_0)_8\,(0) \qquad (18)$$

Expanding

$$h_8(Q_8) = 1 + (h^8)_8(0)Q_8 + 0(Q_8) \equiv 1 + h_0^{'}Q_8 + 0Q_8$$

And substituting it into (18), we arrive at

$$u_8(t,0) \approx \int_0^\infty \beta_8\,(a_8)u_8(t,a_8)da_8 + (h_{a_8}^{'})(\rho_0)_8 Q_8$$

The linearization is the following :

$$[u_8] + [u_8] = -(\mu_8(a_8) + [t_{80}]_8 u_8 - [\alpha_8]_8(a_8)N_8 - [\alpha_8]_8(a_8)Q_8$$

$$u_8(t,0) = \int_0^\infty \beta_8(a_8)u_8(t,a_8)da_8 + [\xi_8]Q_8 \qquad (19)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8$$

where

$$[\alpha_8]_8(a_8) = [t_8^8 a]_8[p,o]_8(a_8)d2(a_8) = [f_{80}^8]_8(\rho\, o\,)_8(a_8)\xi = [\mathrm{h\ o}^8]_8[\rho \mathrm{o}]_8(\mathrm{o}) \text{ and } \mathrm{b}_8 = \mathrm{b}_8 - g_8(P_8)$$

If $g_8(P_0) > b$, clearly, $(\rho_0)_8(a_8), 0)$ is unstable and the predator population is persistent. Assume $g_8(P_0)_8 < b$.

Then, $Q_8(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_8 + u_a = -(\mu_8(a_8) + (f_{80})_8)u_8 - (\alpha_8)_8(a_8)N_8$$

$$= -(\mu_8(a_8) + f_{80})u_8 - \alpha_8(a_8)\int_0^\infty u_8(t,a_8)da_8$$

$$u_8(t,0) = \int_0^\infty \beta_8\,(a_8)u_8(t,a_8)da_8$$

Using the method developed by Cushing [2, 3], we suppose

$$u_8(t,a_8) = \tilde{u}_8(a_8)^{e^{\gamma(t-a_8)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_8}\,da_8$. Then,

A solution of the first differential equation for $\tilde{u}(a_8)$ gives,

$$\tilde{u}_8(a_8) = e^{-\int_0^a \mu_8(\tau_8)d\tau - (f_{80})_8 a}\left(\tilde{u}_8(0) - w\int_0^a e^{\int_0^\tau \mu_8(s_8)ds + (f_{80})}\tau_8\,(\alpha_8)_8(\tau_8)e^{\gamma_8\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\right)\tilde{u}_8(0)$$

$$+w\int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\int_0^a (\alpha_8)_8(\tau_8)e^{-\gamma_8\tau_8}e^{(f_{80})_8\tau_8 + \int_0^\tau \mu_8(s_8)ds}d\tau_8\,da_8$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}da\right)(\tilde{u}_8(0))_8$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_8 a_8-(f_{80})_8 a_8-\int_0^a \mu_8(\tau_8)d\tau}\int_0^a e^{(f_{80})_8\tau_8+\int_0^\tau \mu_8(s_8)ds+y\tau}\alpha_8(\tau)d\tau\, da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_8 a_8-(f_{80})_8 a_8-\int_0^a \mu_8(t_8)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_8\,(a_8)Y_8(\gamma_8,a_8)da_8\right)u_8(0)$$

$$+w\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)\int_0^a(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8=0,$$

(24)

$$\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\tilde{u}_8(0)$$

$$+\left(-1-\int_0^{\infty}Y_8(\gamma_8,a_8)\int_0^a(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8\right)w=0.$$

Equation (20) has a solution of the form of $u_8(t,a_8)=u_8(a_8)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_8(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_8(\gamma_8,a_8)=\int_0^a(\rho_0)_8(\tau_8)Y^{-8}\,(\gamma_8,a_8)d\tau_8$. We obtain the characteristic equation

$$f_{80_8}^{l}\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\left(\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$+\left(1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)da\right)\left(1+(f_{80}^{l})_8\int_0^{\infty}Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_8(a_8) = (\rho_0)_8(0)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - (f_{80})_8 a_8}$$

$$I_8(\gamma_8, a_8) = \int_0^a (\rho_0)_8(\tau_8)e^{\gamma_8\tau_8 + (f_{80})_8\tau + \int_0^\tau \mu_8(s_8)ds_8}d\tau$$

$$= \frac{(\rho_0)_8(0)}{\gamma_8}(e^{\gamma_8 a_8} - 1)$$

Let

$$E_8(a_8) = e^{-(f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau_8}$$

Then $[Y_8]_8[\gamma_8]_8, a_8) = E_8(a_8)e^{-\gamma_8 a_8}$. by a simple calculation,

$$\int_0^\infty \beta_8(a_8)Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma}\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}\ da\right);$$

$$\int_0^\infty Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma_8}\left(\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}\ da_8\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_8(0)}{\gamma_8}[f_{80}^8]\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}da_8\ (\int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8)$$

$$+\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right).$$

$$\left(1+\frac{(\rho_0)_8(0)}{\gamma_8}(f_{80}^{'})_8\left(\int_0^{\infty}E_8(a_8)da_8-\int_0^{\infty}E_8(a_8)e^{-\gamma_8 a_8}da_8\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right)(1+\frac{(f_{80}^{'})_8}{\gamma_8}(\rho_0)_8(0)\int_0^{\infty}E_8(a_8)da_8)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)da_8=1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\neq 0,$$

Since $\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)\int_0^{a}\alpha_8(\tau_8)Y^{-8}(\gamma_8,\tau_8)d\tau_8 da_8>0$. Therefore, the characteristic solution is

$$\gamma_8=-(f_{80}^{'})_8(\rho_0)_8\int_0^{\infty}E_8(a_8)da_8.$$

Since $f_8^{'}(P_0)_8>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_8),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_8),0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_8 + (\mu_8(a_8) + [f_8]_8(P_8^*) + [f_8]_8(Q_8^*))P_8^* = 0, \quad (26)$$

$$p_8^*(0) = h_8(Q_8^*)\int_0^\infty \beta_8(a_8)p_8^*(a_J)p_8^*(a_8)da_8 \quad (27)$$

$$g_8(P_8^*) - b_8 = 0, \quad (28)$$

Where $P_8^* = \int_0^\infty p_8^*(a_8)\, da_8$

Equation (28) yields $P^* = g^{-8}(b)$,

From (26),

$$p_8^*(a) = p_8^*(0)e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}$$

Substituting into (27), we need to solve

$$\left(1 - h_8(Q_8^*)\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}da_8\right)p_8^*(0) = 0$$

For $Q_8^*$

Fix $p_8^*(0)$ and $P_8^*$ Let

$$H_8(z_8) = 1 - h_8(z_8)\int_0^\infty \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}da_{,8}$$

And assume that

$$[k_3]_8 \int_0^{\infty} \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(g_8^{-8}(b_8))a_8 - [k_8]_8 a_8}\, da_8 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty} [H_8(z_8)] > 0.$$

Clearly,

$$\frac{dH_8}{dz_8} = h_8(z)(\,[f_8^{'}]_8(z_8)\,) \int_0^{\infty} a_8 \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - f_8(P_8^*)a_8 - f_8(z_8)a_8} da_8$$

$-h^{'}(z_8) \int_0^{\infty} \beta_8(a_8)\, e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}\; da_8$

By the assumptions on $\left[f_8^8\right]_8 and\ [h^8]^8 dH_8/dz > 0$. Hence, if

$$H_8(0) = 1 - \int_0^{\infty} \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8} da_8 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_8(0) < 0$ follows if $g_8^-(b_8) < [f_8^{-8}]_8(r_8)$, where $r_8$ is the solution of (12).

Indeed, since

$$R_8(s_8) = \int_0^{\infty} \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 e^{-y_8^a}} da_8$$

Is a decreasing function of $s_8$ when $P_8^* = [g^{-8}]_8(b_8) < [f_8^{-8}]_8(r_8)$, that is $[f_8]_8(P_8^*) < r$, $R\big(f_8(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_8(g^{-8}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8\left(g_8^{-8}(b)\right)a_8} > 1. \qquad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_8^*(a_8), Q_8^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_8^*(a_8), Q_8^*)$ is more complicated. We again consider perturbations $m_8(t_8 a_8) = p_8(t_8, a_8) - p_8^*(a_8)$ and $W_8(t) = Q_8(t) - Q^*$.Then

$$M_8(t) \equiv \int_0^\infty m_8(t, a_8)da_8 = P_8(t) - P_8^*,$$

$$[m_8]_8 + [m_a]_8 = [p_8]_8 + [p_a]_8 - \left[p^*\right]_8(a_8) \approx -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8$$

$$-[f_{8*}^8]_8 p_8^*(a_8)M - [f_{8*}^8]_8 p_8^*(\mathrm{a})W_8,$$

$$m_8(t,0)p_8(t,0) - p_8^*(0) \approx h_* \int_0^\infty \beta_8(a_8)m_8(t,a_8)\,da + \frac{[h_*^8]_8}{[h_*]_8}\, p_8^*(0)W_8$$

$$\frac{dW_8}{dt} = \frac{dQ_8}{dt} \approx [g_*^8]_8 Q_8^* M_{,8}$$

Where $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $\left[f_{8*}^{'}\right]_8$ denoting $\left[f^{'}{}_8(P^*)\right]_8\left[f_{8*}^{'}\right]_8$ denoting $[f_8]_8$ $(Q_8^*,$ $h_{*}^{'}$ denoting $h_*^{'}$ denoting $h_8^{'}$ $(Q_8^*)$, and $\left[g_*^{'}\right]_8$

Denoting $\left[\dot{g}\right]_8(P_8^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_8]_8 + [m_a]_8 = -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8 - \left[\dot{f}_{8*}\right]_8[p_8^*](a_8)M_8 - \left[\dot{f}_{8*}\right]_8 p_8^*(a_8)W_8$$

$$m_8(t,0) = [h_*]_8 \int_0^\infty \beta_8(a_8)m_8(t,a_8)da_8 + \frac{[h_*^8]_8}{[h^*]_8}[p^*]_8(0)\,W_{,8} \qquad (30)$$

$$\frac{dW_8}{dt} = [g_*^8]_8[Q_8^*]M_8$$

Integrating $[m_8]_8(t,a_8)$ with respect to a from 0 to ∞, we have

$$\frac{dM_8(t)}{dt} = \int_0^\infty [mt]_8(t,a_8)da_8$$

$$= -\int_0^\infty [m_a]_8\,(t,a_8)da_8 - \int_0^\infty \mu(a_8)m(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + \left[\dot{f}_{8*}\right]_8 \int_0^\infty p_8^*(a_8)da_8\right)M_8 - [f_{8*}^8]_8 \int_0^\infty p_8^*(a_8)da_8 W_8$$

$$= m_8(t,0) - \int_0^\infty \mu_8(a_8)\, m_8(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + \left[\dot{f}_{8*}\right]_8 P_8^*\right)M_8 - \left[\dot{f}_{8*}\right]_8 P_8^* W_{,8}$$

From (30), we arrive at the equations

$$\frac{dW_8(t)}{dt} = \int_0^\infty ([h_*]_8\beta_8(a_8) - \mu_8(a_8))m_8(t,a_8)da_8$$

$$-\left([f_{8*}] + [f_{8*}]_8 + [f_{8*}']_8 P_8^*\right) M_8 + \left(\frac{[h_*']_8}{[h_*]^8}\, P_8^*(0) - f_{8*}' P_8^*\right) W_8,$$

(31)

$$\frac{dW_8(t)}{dt} [g'_*]_8\, Q_8^*\, M_{.8}$$

By the mean value theorem,

$$\int_0^\infty ([h_{"*}]_8 \beta_8(a_8) - \mu_8(a_8)) m_8(t, a_8) da_8 = ([h_*]_8 \beta_8(\theta_8) - \mu_8(\theta_8)) M_{,8}$$

Where $0< \theta_8 = \theta_8(t) < \infty$.

And if a $^8_8$ [a $^8_8$]2<4 [a $_8$]2<4 a $_8$e,

Re $\lambda_8 [\lambda_8]_8 = \text{Re}\lambda_8 [\lambda_8]_8 = -\frac{8}{8} a_8$[a2 ] $_8 < 0$.

Let ⊿=max$_{8<u<\infty}$ |h $_*$ [h $_*$] $_8$β $_8$($a_8$) − $\mu_8(a_8)$|. Then|$n_8$(t)| ⊿∀t . Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium(ρ $^*$ (a), $\boldsymbol{Q}^*$) is asymptotically stable if either a $^8_8$ <4 a $_8$ e and [a $_8$] > 1⊿,or [a $^8_8$] [a $_8$]e and (a $_8$ − $\sqrt{(a^8_8)_8 - 4(a_8)_8 e}$ > 1⊿

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+8}\right)$ and β $_8$($a_8$) = ae-$^{ma}$ Then ⊿=[$\mu_0$] + 2 Simple calculation leads to

$$\int_0^{\infty} \beta_8(a_8)\mathrm{e}^{-\int_0^{\mathrm{a}} \mu_8(\tau_8)+d\tau_8}\mathrm{da} = \frac{1}{(\mathrm{m}+\mu_0+1)^8}+\frac{1}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_8 + [\mu_0]$ 0.5. Assume$[f_8](x_8)= K_8x_8$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_8(x_8)= [0x_8]/(2 + x_8)$; the functional responses $[[f_8](x) = [0]\, x_8/(1 + x_8)$and $h_8(x_8)$=$(1+\omega_8x_8)/(1 + x_8)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_8 + [\mu_0] + [\mathrm{K}][P_0])}{[m_8]_8 + (\mu_0)_8 + [\mathrm{K}][P_0])_8^3\ [\mathrm{K}]_8[P_0]_8^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_8$i e $[P_0] < b/(\delta_8 - b_8)$then the equilibrium$(P_0)$ $[P_0]_8(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_8$

$[f_{8*}] = K_8\ \mathrm{P}^*\ [f_{8*}] = \frac{\sigma \mathrm{Q}_8^*}{8+\mathrm{Q}_8^*}$ $[h_*] = \frac{8+\omega_8\mathrm{Q}_8^*}{8+\mathrm{Q}_8^*}$

$[f_{8*}] = K_8\ [\mathrm{f}_{8*}^8] = \frac{\sigma_8}{(8+\mathrm{Q}_8^*)^8}[\ \mathrm{g}_*^8] = \frac{\sigma_8}{(8+\mathrm{P}_8^*)^8}$ $[\mathrm{h}_*^8] = \frac{\omega_8-8}{(8+\mathrm{Q}_8^*)^8}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_8}{(m+[\mu_0]_8+1+\frac{K_8b_8}{\delta_8+b_8}+\sigma)^8}+\frac{1\omega_8}{([m_8]_8+[\mu_0]_8+1+\frac{K_8b_8}{\delta_8+b_8}+\sigma_8)^3}<1$$

And

$$\frac{1}{(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^8} + \frac{1\omega_8}{(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8} + (m_8 + [\mu_0]_8 + \frac{K_8 b_8}{\delta_8 - b_8} + \sigma_8)\, Q_8^*)^3$

$= m_8 + [\mu_0]_8 + 3 + \frac{K_8 b_8}{\delta_8 - b_8}(1 + Q^*)(1 + \omega_8 Q_8^*$

Now

$$\left[a_{11}\right] - 4\left[a_{8}\right] e = \left(\frac{1K_8 b_8}{\delta_8 - b_8}\right)^8 + \frac{4K_8 b_8 \sigma_8 Q_8^*}{(\delta_8 - b_8)(8 + Q_8^*)} + \frac{\sigma_8^8 Q_8^{*8}}{(1 + Q_8^*)^8}$$

$$\frac{4\, b_8 \sigma_8 (\delta_8 - b) Q^*}{\delta_8 (1 + Q_8^*)^8} - \frac{4\, (\delta_8 - b_8)^8 (1 - \omega_8) Q_8^*}{\delta_8 (1 + \omega_8 Q_8^*)(1 + Q_8^*)} P^*(0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_8\, ae_8^{mu}\ [\mu_0]_8 - \frac{a_8}{a_8 + 8}$$

Hence if $[a_8^8 > 4(a_8)_8 e$ and

$$[a_8] = \frac{1K_8 b_8}{\delta_8 - b_8} + \frac{\sigma_8 Q_8^*}{1 + Q_8^*} > 1\Delta$$

**The nontrivial equilibrium is stable.**

IF we define

$$[a_8]_8 = [f_{8*}]_8 + [f_{8*}]_8 + [f_{8*}] P_8^*$$

$$[a_8]_8 = \left[f_{8*}^{8}\right] P_8^* - \frac{h_*^8}{h_*} P^*(0)$$

$$e = [g_*^8]_8 Q_8^*$$

$$\eta_8(t) = \left[h_*\right]_8 \beta_8 \left(\theta_8(t)\right) - \mu_8(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\left[a_8\right] M_8 - [a_8]_8 W_8 - [n(t)] M_8$$

$$w_8(t) = eM_8$$

Write

$$A = \begin{pmatrix} -[a_8]_8 - [a_8]_8 \\ e \qquad 0 \end{pmatrix}$$

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| $B=(t)=\begin{pmatrix} [\eta(t)]_8 & 0 \\ 0 & 0 \end{pmatrix}$ |
| And |
| $X(t)=\begin{pmatrix} M_8 \\ W_8 \end{pmatrix}$ |
| The solution of (32)is |
| $X_8(t) = e^{A,t}[X]_8(0) + \int_0 e^{At}$ |
| This equation notwithstanding its importance in all the models stand out for its importance for the following models<br><br>to which it is applied: |
| **Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016** |

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| (323)These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field. |
| (324)The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets. |
| (325) Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state. |
| (326) Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands. |
| (327)The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI |
| (328) The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9]. |
| (329)This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive. |

(330)More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

(331)Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.

(332)Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.

(333) The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.

(334)This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.

(335)The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.

(336)This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.

(337) For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.

(338)The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(339)Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.

(340)We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.

(341)When this is true—as is often the case—the system serves as (=) a "bath" for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called **thermalization.**

(342)However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.

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| (343)Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1). |
| (344)Bringing them to thermal equilibrium requires (e) an external bath. |
| (345) But so far, there has been no rigorous proof of (e) the existence of **such many-body localization**. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption. |
| (346)The concept of **localization** was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations). |
| (347)The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century. |
| (348)But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information. |
| (349)Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment. |
| (350)For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood. |
| (351) It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7]. |
| (352) Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature. |
| (353) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators. |
| (354)However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8]. |
| (355)Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems. |
| (356) Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL. |
| (357)As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked. |
| (358)His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice. |
| (359) An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site. |

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| (360)This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself). |
| (361)This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition. |
| (362)However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime: |
| (363) These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects. |
| (364)It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies). |
| (365) This is a realistic scenario for (e) all known systems. |
| (366)On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects. |
| (367)On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other. |
| (368)No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true. |
| (369) His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found. |
| (370)Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization? |
| (371)Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects? |
| (372)Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions. |
| (373)So in those systems, the quantum phase transition does indeed exist. |
| (374)But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension. |
| (375)Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters. |

# PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION ATTRITION MODELS

## MODULE NINE

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Model is applied to following formulations Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society**

**Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice ($\alpha$=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP This paper is divided into four parts. In part A, some general considerations about (e&eb) gravitational radiation are followed by (e) a treatment of the scalar wave equation in the manner later to be applied to (e&eb) Einstein's field equations. In part B, a co-ordinate system is specified which is suitable for (e) investigation of outgoing gravitational waves from (e) an isolated axi-symmetric reflexion-symmetric system. The metric is expanded in (eb) negative powers of (e) a suitably defined radial co-ordinate r, and the vacuum field equations are investigated in detail. It is shown that the flow of information to infinity is controlled by (e&eb) a single function of (e) two variables called the news function. Together with initial conditions specified on a light cone, this function fully defines (eb) the behaviour of the system. No constraints of (e) any kind are encountered. In part C, the transformations leaving (eb) the metric in the chosen form are determined. An investigation of the corresponding transformations in (eb) Minkowski space suggests (eb) that no generality is lost by (e) assuming that the transformations, like the metric, may be expanded in (eb) negative powers of r. Gravitational Waves in General Relativity VII Waves from Axi-Symmetric Isolated Systems Royal Society Proceedings H. Bondi, M. G. J. van der Burg, A. W. K. Metzner Published 21 August 1962.DOI: 10.1098/rspa.1962.0161 STOP There is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant $\Lambda$, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive $\Lambda$ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the $\Lambda>0$ generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible. Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of (e) GW signals from BBHs with component spins $\chi_{1,2}\in[-0.99,0.99]$ aligned (e&eb) with the orbital angular momentum, component masses $m_{1,2}\in[2,48]M_\odot$, and (e&eb) total mass $M_{total}\leq 50\ M_\odot$. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin**

**bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins >+0.6 and component masses ≳20 M⊙. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201**

## 45. MODEL FORMULATION

Let $p_9(t, a_9)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(g_9)Q_9$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_9(0) = 0, \quad g_9(P_9) > 0, \ P_9 \ 0, \qquad \lim_{P_9 \to \infty} g_9(P_9) = (K_9)_9 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_9 + (\rho_a)_9 + D_9(a_9, P_9, Q_9)\rho_9 + 0, \qquad 0 < a_9 \ \infty,$$

$$\rho_9(t,0) = \int_0^\infty n_9(a_9, P_9, Q_9)\rho_9(t_9, a_9)da_9$$

Where $D_9$ is the mortality function and $n_9$ is the fertility function.

Put $_9$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_9}{dt}\right]_9 = x_9\left[\frac{dQ_9}{dt}\right]_9 - x_9\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_9}{dt}\right]_9 = x_3\left[\frac{dQ_3}{dt}\right]_9 - x_4\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_9}{dt}\right]_9 - x_6\left[\frac{dQ_3}{dt}\right]_9$$

Where $x_9, x_9, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_9$ and $g_9(P_9)Q_9$

Now when you add the term $(b_9Q_9)$ with $g_9(P_9)(Q_9)$ we are infect see is that $(b_9Q_9)$ is credited incrementally part of $g_9(P_9)(Q_9)$ or instant nearly, when subtracted from $\frac{dQ_9}{dt}$ , it means that $g_9(P_9)(Q_9)$ is slowly eaten by $\frac{dQ_9}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_9 = \mu_9(a_9) + (f_9)_9(P_9)_9 + (f_9)_9(Q)_9 \quad (2)$$

$$n_9 = \beta_9(a_9)h_9(Q_9), \quad (3)$$

Where $\mu_9$the age- dependent debit rate whit is $\mu_9(a_9)$ 0 $\forall a > 0$; $(f_9)_9$ is the density- dependent death rate with $f_9(0) = 0$, $(f_9^9)_9(P_9)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_9)_9(P_9) = (k_9)_9$ ; $f_9$ is the functional response of creditors to prey with $[f_9]_9(0) = 0, (f_9^9)_9(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_9)_9(Q_9) = [k_9]$ ; $\beta$ is the age dependent birth

rate with $\beta(a) \geq 0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h'(Q) \leq 0, \forall Q \geq 0$, and $\lim_{Q\to\infty} h(Q) = k_3 \geq 0$.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_9 + (\rho_a)_9 = -\mu_9(a_9) + (f_9)_9 P_9) + \big(f_9)_9(Q_9)\big)(\rho_9)_9 \tag{4}$$

$$\rho_9(t_9, 0) = h_9(Q_9) \int_0^\infty \beta_9(a_9)\rho_9(t, a_9)da, \tag{5}$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(P_9)Q_9 \tag{6}$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_9$ and $Q_9$ of (4-6) satisfy

$$0 \leq \rho_9(t, a_9) \leq \rho_9^*(t, a_9),$$

$$0 \leq Q_9(t) \leq Q_9^*(t),$$

Where $\rho_9^*(t, a_9)$ and $Q_9^*(t)$, are solutions of

$$(\rho_t^*)_9 + (\rho_a^*)_9 = -(\mu_9(a_9) + (f_9)_9(P_9^*) + (f_9)_9 Q_9^*))\, \rho_9^*, \tag{7}$$

$$\rho_9^*(t, 0) = \int_0^\infty \beta_9(a_9)\, \rho_9^*(t, a_9)da_9 \tag{8}$$

$$\frac{dQ_9^*}{dt} = -b_9 Q_9^* + g_9(P_9^*)Q_9^*, \tag{9}$$

With $P_9^*(t) = \int_0^\infty \rho_9^*\ (t, a_9)da_9$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_9^*(t)}{dt} = P_9^*(t)\big(r_9 - (f_9)(P_9^*) - (f_9)_9(Q^*)\big), \quad (10)$$

$$\frac{dQ_9^*(t)}{dt} = Q_9^*(t)(-b_9 + g_9(P_9^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_9(a)e^{-\int_0^a \mu_9(\tau)dt}e^{-ra}\ da_9 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_9^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(t)dt}\ da_9 > 1, \quad (H1)$$

Which guarantees $r_9 > 0$.

*Remark.* If $\beta_9(a) = [\beta_0]_9 a_9 e^{-\alpha_9 a_9}, \mu_9(a_9) \equiv \mu_9$ then the solution of (12) is

$$r_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_9]_9(P_9) = 0, [f_9](Q_9) = \lambda_9 Q_9\ h_9(Q_9) = 1, \quad and\ \ g_9(P_9) = c_9 P_9$$

$$[\rho_9]_9 + [\rho_a]_9 = -(\mu_9 + \lambda_9 Q_9)\rho_9$$

$$\rho_9(t,0) = \int_0^{\infty} [\beta_0]_9 a_9 e^{-\alpha_9 a_9} \rho_9(t,a_9)\, da_9$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + c_9 P_9 Q_9$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_9}{dt} = \omega_9 P_9 - \lambda_9 P_9 Q_9$$

$$\frac{dQ_9}{dt} = b_9 Q_9 - c_9 P_9 Q_9$$

Where $\omega_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$

THE EQUILIBRIUM $(\rho_0)_9(a_9), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_9(a_9), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_9 + (\rho_a)_9 = -\big(\mu_9(a_9) + (f_9)_9(P_9)\big)\rho,$$

$$\rho_9(t,0) = \int_0^{\infty} \beta_9\,(a_9)\rho_9(t,a_9)da_9 \qquad (13)$$

The solution of (13) has the form of

$$\rho_9(t,a_9) = Ce_9^{r(t-a_9)} - \int_0^a (\tau)d\tau - \int_0^9 (f_9)_9\big(P_9(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_9(t) = C_9 e^{r_9 t - \int_0^t (f_9)_9(P_9(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_9^{-9}(r)$ if $k_9 > r$. This is equivalent to

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{(k_9)_9 a_9} da_9 < 1.$$

In fact, if we define

$$R_9(s_9) = \int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau)d\tau} e^{-s_9 a_9}\, da_9$$

$R(s)$ is a decreasing function. $R_9(r_9) = 1$ And $R_9(k_9)_9 < 1$ imply $r_9 < (k_9)_9$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_9(a_9), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{k_9 a_9} da_9 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_9 \equiv \int_0^\infty (\rho_0)_{99}(a_9)da_9$ *satisfies*

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{-(f_9)_9(P_0)a_9} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_9(a_9), 0)$

Let $u_9(t,a_9) = \rho(t,a_9) - \rho_0(a_9)$ Then

$$N_9(t) \equiv \int_0^\infty u_9\,(t, a_9)da_9 = \int_0^\infty \rho_9(t, a_9)da_9 - \int_0^\infty (\rho_0)_9(a_9)da = P_9(t) - P_0)_9$$

And

$$(u_9)_9 + (u_a)_9 = (\rho_t)_9 + (\rho_a)_9 - (\rho_0^9)_9(a_9)$$

$$= -\mu_9(a_9) + (f_9)_9(P_9) + (f_9)_9(Q_9))\,(u_9 + (\rho_0)_9) + \mu_9(a_9)(\rho_0)_9 + (f_9)_9(P_0)_9(\rho_0)_9$$

$$= -\mu_9(a_9)u_9 - (f_9)_9(P_9) + (f_9)_9(Q_9))\,u_9 - (\rho_0)_9\,(f_9)_9(P_9) - (f_9)_9(P_0)_9$$

$$+ (f_9)_9(Q_9)), \qquad (15)$$

Where $(\rho_0^9)_9(a_9)$ denotes $d(\rho_0)_9(a_9)/da_9$

Expansion $(f_9)_9(P_9)$ and $(f_9)_9(Q_9)$ at $(P_0)_9$ and 0 yield respectively

$$(f_9)_9(P_9) = (f_9)_9((P_0)_9 + N_9)$$

$$= (f_9)_9(P_0)_9 + (f_9^9)_9(P_0)_9N_9 + o_9(N_9), \qquad (16)$$

$$(f_9)_9(Q_9)$$

$$= (f_9^9)_9(0)Q + o_9(Q_9) \qquad (17)$$

Write $f_9(P_0) = f_{90}, f_9'(P_0) = f_{90}'$, and $f_9'(0) = f_{90}'$ and substitute (16) and (17) into (15) to obtain

$$(u_9)_9 + (u_a)_9 \approx -(\mu_9(a_9) + (f_{90})_9)u - (f_{90}^9)_9(\rho_0)_9(N_9) - (f_{90}^9)_9(\rho_0)_9(Q_9)$$

$$u_9(t, 0) = \rho_9(t, 0)(\rho_0)_9(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t, a) + (\rho_0)_9(a_9)da_9 - (\rho_0)_9\,(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t, a_9)da_9 + (h_9(Q_9) - 1)(\rho_0)_9\,(0) \qquad (18)$$

Expanding

$$h_9(Q_9) = 1 + (h^9)_9(0)Q_9 + 0(Q_9) \equiv 1 + h_0'Q_9 + 0Q_9$$

And substituting it into (18), we arrive at

$$u_9(t,0) \approx \int_0^\infty \beta_9(a_9)u_9(t,a_9)da_9 + (h_{a_9}')(\rho_0)_9Q_9$$

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PREDATOR-PREY POPULATION MODELS

The linearization is the following :

$$[u_9] + [u_9] = -(\mu_9(a_9) + [t_{90}]_9u_9 - [\alpha_9]_9(a_9)N_9 - [\alpha_9]_9(a_9)Q_9$$

$$u_9(t,0) = \int_0^\infty \beta_9(a_9)u_9(t,a_9)da_9 + [\xi_9]Q_9 \qquad (19)$$

$$\frac{dQ_9}{dt} = -b_9Q_9$$

where

$$[\alpha_9]_9(a_9) = [t_9^9a]_9[p,o]_9(a_9)d2(a_9) = [f_{90}^9]_9(\rho\, o\,)_9(a_9)\xi = [\mathrm{h\ o}^9]_9[\rho\mathrm{o}]_9(\mathrm{o}) \text{ and } \mathrm{b}_9 = \mathrm{b}_9 - g_9(P_9)$$

If $g_9(P_0) > b$, clearly, $(\rho_0)_9(a_9), 0)$ is unstable and the predator population is persistent. Assume $g_9(P_0)_9 < b$.

Then, $Q_9(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_9 + u_a = -(\mu_9(a_9) + (f_{90})_9)u_9 - (\alpha_9)_9(a_9)N_9$$

$$= -(\mu_9(a_9) + f_{90})u_9 - \alpha_9(a_9)\int_0^\infty u_9(t,a_9)da_9$$

$$u_9(t,0) = \int_0^\infty \beta_9\,(a_9)u_9(t,a_9)da_9$$

Using the method developed by Cushing [2, 3], we suppose

$$u_9(t,a_9) = \tilde{u}_9(a_9)^{e^{\gamma(t-a_9)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_9}\,da_9$. Then,

A solution of the first differential equation for $\tilde{u}(a_9)$ gives,

$$\tilde{u}_9(a_9) = e^{-\int_0^a \mu_9(\tau_9)d\tau-(f_{90})_9 a}\left(\tilde{u}_9(0) - w\int_0^a e^{\int_0^\tau \mu_9(s_9)ds+(f_{90})}\tau_9\,(\alpha_9)_9(\tau_9)e^{\gamma_9\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\right)\tilde{u}_9(0)$$

$$+w\int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\int_0^a (\alpha_9)_9(\tau_9)e^{-\gamma_9\tau_9}e^{(f_{90})_9\tau_9+\int_0^\tau \mu_9(s_9)ds}d\tau_9\,da_9$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}da\right)(\tilde{u}_9(0))_9$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\int_0^a e^{(f_{90})_9\tau_9+\int_0^\tau \mu_9(s_9)ds+y\tau}\alpha_9(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(t_9)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_9\,(a_9)Y_9(\gamma_9,a_9)da_9\right)u_9(0)$$

$$+w\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)\int_0^a(\alpha_9)_9(\tau)_9(Y^{-9})_9(\gamma_9,\tau_9)d\tau_9\,da_9=0,$$

(24)

$$\left(\int_0^{\infty}Y_9(\gamma_9,a_9)da_9\right)\tilde{u}_9(0)$$

$$+\left(-1-\int_0^{\infty}Y_9(\gamma_9,a_9)\int_0^a(\alpha_9)_9(\tau)_9(Y^{-9})_9(\gamma_9,\tau_9)d\tau_9\,da_9\right)w=0.$$

Equation (20) has a solution of the form of $u_9(t,a_9)=u_9(a_9)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_9(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_9(\gamma_9,a_9)=\int_0^a(\rho_0)_9(\tau_9)Y^{-9}\,(\gamma_9,a_9)d\tau_9$. We obtain the characteristic equation

$$f_{90_9}^{l}\left(\int_0^{\infty}Y_9(\gamma_9,a_9)da_9\right)\left(\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)I_9(\gamma_9,a_9)da_9\right)$$

$$+\left(1-\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)da\right)\left(1+(f_{90}^{l})_9\int_0^{\infty}Y_9(\gamma_9,a_9)I_9(\gamma_9,a_9)da_9\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_9(a_9) = (\rho_0)_9(0)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - (f_{90})_9 a_9}$$

$$I_9(\gamma_9, a_9) = \int_0^a (\rho_0)_9(\tau_9)e^{\gamma_9\tau_9 + (f_{90})_9\tau + \int_0^\tau \mu_9(s_9)ds_9}d\tau$$

$$= \frac{(\rho_0)_9(0)}{\gamma_9}(e^{\gamma_9 a_9} - 1)$$

Let

$$E_9(a_9) = e^{-(f_{90})_9 a_9 - \int_0^a \mu_9(\tau_9)d\tau_9}$$

Then $[Y_9]_9[\gamma_9]_9, a_9) = E_9(a_9)e^{-\gamma_9 a_9}$. by a simple calculation,

$$\int_0^\infty \beta_9(a_9)Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma}\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}\ da\right);$$

$$\int_0^\infty Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma_9}\left(\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}\ da_9\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_9(0)}{\gamma_9}[f_{90}^9]\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}da_9\ (\int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9)$$

$$+\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right).$$

$$\left(1+\frac{(\rho_0)_9(0)}{\gamma_9}(f_{90}^{'})_9\left(\int_0^{\infty}E_9(a_9)da_9-\int_0^{\infty}E_9(a_9)e^{-\gamma_9 a_9}da_9\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right)(1+\frac{(f_{90}^{'})_9}{\gamma_9}(\rho_0)_9(0)\int_0^{\infty}E_9(a_9)da_9)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (94), to

$$1-\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)da_9=1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\neq 0,$$

Since $\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)\int_0^{a}\alpha_9(\tau_9)Y^{-9}(\gamma_9,\tau_9)d\tau_9 da_9>0$. Therefore, the characteristic solution is

$$\gamma_9=-(f_{90}^{'})_9(\rho_0)_9\int_0^{\infty}E_9(a_9)da_9.$$

Since $f_9^{'}(P_0)_9>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_9),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_9),0)$ which is related to the extinction of predators (debit summation types: that which acts as a resource)and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_9 + (\mu_9(a_9) + [f_9]_9(P_9^*) + [f_9]_9(Q_9^*))P_9^* = 0, \quad (26)$$

$$p_9^*(0) = h_9(Q_9^*)\int_0^\infty \beta_9(a_9)p_9^*(a_J)p_9^*(a_9)da_9 \quad (27)$$

$$g_9(P_9^*) - b_9 = 0, \quad (28)$$

Where $P_9^* = \int_0^\infty p_9^*(a_9)\, da_9$

Equation (28) yields $P^* = g^{-9}(b)$,

From (26),

$$p_9^*(a) = p_9^*(0)e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}$$

Substituting into (27), we need to solve

$$\left(1 - h_9(Q_9^*)\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}da_9\right)p_9^*(0) = 0$$

For $Q_9^*$

Fix $p_9^*(0)$ and $P_9^*$ Let

$$H_9(z_9) = 1 - h_9(z_9)\int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9}\, da_{,9}$$

And assume that

$$[k_3]_9 \int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(g_9^{-9}(b_9))a_9 - [k_9]_9 a_9}\, da_9 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_9(z_9)] > 0.$$

Clearly,

$$\frac{dH_9}{dz_9} = h_9(z)\left( [f_9^{'}]_9(z_9) \right)\int_0^\infty a_9\beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - f_9(P_9^*)a_9 - f_9(z_9)a_9} da_{\square}$$

$-h^{'}(z_9)\int_0^\infty \beta_9(a_9)\, e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9}\quad da_9$

y the assumptions on $\left[f_9^9\right]_9 and\ [h^9]^9 dH_9/dz > 0$. Hence, if

$$H_9(0) = 1 - \int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9} da_9 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_9(0) < 0$ follows if $g_9^-(b_9) < [f_9^{-9}]_9(r_9)$, where $r_9$ is the solution of (12).

Indeed, since

$$R_9(s_9) = \int_0^{\infty} \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 e^{-y\frac{a}{9}} da_9}$$

Is a decreasing function of $s_9$ when $P_9^* = [g^{-9}]_9(b_9) < [f_9^{-9}]_9(r_9)$, that is $[f_9]_9(P_9^*) < r$, $R\big(f_9(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_9(g^{-9}(b))) > 1, i.e.,$

$\int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9\left(g_9^{-9}(b)\right)a_9} > 1.$ (H4)

Then the system (4)-(6) has a unique positive equilibrium $(P_9^*(a_9), Q_9^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_9^*(a_9), Q_9^*)$ is more complicated. We again consider perturbations $m_9(t_9 a_9) = p_9(t_9, a_9) - p_9^*(a_9)$ and $W_9(t) = Q_9(t) - Q^*$.Then

$M_9(t) \equiv \int_0^{\infty} m_9(t, a_9) da_9 = P_9(t) - P_9^*,$

$$[m_9]_9 + [m_a]_9 = [p_9]_9 + [p_a]_9 - \left[p^*\right]_9(a_9) \approx -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9$$

$$-[f_{9*}^9]_9 p_9^*(a_9)M - [f_{9*}^9]_9 p_9^*(\mathrm{a})W_9,$$

$$m_9(t,0)p_9(t,0) - p_9^*(0) \approx h_* \int_0^{\infty} \beta_9(a_9) m_9(t, a_9)\, da + \frac{[h_*^9]_9}{[h_*]_9}\, p_9^*(0)W_9$$

$$\frac{dW_9}{dt} = \frac{dQ_9}{dt} \approx [g_*^9]_9 Q_9^* M_{,9}$$

Where $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $\left[f_{9*}^{'}\right]_9$ denoting $\left[f^{'}{}_9(P^*)\right]_9 \left[f_{9*}^{'}\right]_9$ denoting

$[f_9]_9$ $(Q_9^*,$ $h_9^*$ denoting $h_*^{'}$ denoting $h_9^{'}$ $(Q_9^*)$, and $\left[g_*^{'}\right]_9$

Denoting $\left[g^{'}\right]_9(P_9^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_9]_9 + [m_a]_9 = -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9 - \left[f^{'}{}_{9*}\right]_9 [p_9^*](a_9)M_9 - \left[f_{9*}^{'}\right]_9 p_9^*(a_9)W_9$$

$$m_9(t,0) = [h_*]_9 \int_0^\infty \beta_9(a_9)m_9(t,a_9)da_9 + \frac{[h_*^9]_9}{[h^*]_9}[p^*]_9(0)\, W_{,9} \qquad (30)$$

$$\frac{dW_9}{dt} = [g_*^9]_9[Q_9^*]M_9$$

Integrating $[m_9]_9(t,a_9)$ with respect to a from 0 to ∞, we have

$$\frac{dM_9(t)}{dt} = \int_0^\infty [mt]_9(t,a_9)da_9$$

$$= -\int_0^\infty [m_a]_9\,(t,a_9)da_9 - \int_0^\infty \mu(a_9)m(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + \left[f^{'}{}_{9*}\right]_9 \int_0^\infty p_9^*(a_9)da_9\right)M_9 - [f_{9*}^9]_9 \int_0^\infty p_9^*(a_9)da_9 W_9$$

$$= m_9(t,0) - \int_0^\infty \mu_9(a_9)\, m_9(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + \left[f_{9*}^{'}\right]_9 P_9^*\right)M_9 - \left[f_{9*}^{'}\right]_9 P_9^* W_{.9}$$

From (30), we arrive at the equations

$$\frac{dW_9(t)}{dt} = \int_0^{\infty} ([h_*]_9\beta_9(a_9) - \mu_9(a_9))m_9(t,a_9)da_9$$

$$-\left([f_{9*}] + [f_{9*}]_9 + \left[f_{9*}^{'}\right]_9 P_9^*\right) M_9 + \left(\frac{\left[h_*^{'}\right]_9}{[h_*]^9}\ P_9^*(0) - f_{9*}^{'} P_9^*\right) W_9,$$

(31)

$$\frac{dW_9(t)}{dt}\left[g^{'}{}_*\right]_9 Q_9^*\, M_{.9}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{"*}]_9\beta_9(a_9) - \mu_9(a_9))m_9(t,a_9)da_9 = ([h_*]_9\beta_9(\theta_9) - \mu_9(\theta_9))M_{,9}$$

Where $0< \ \theta_9 = \theta_9(t) < \infty$.

And if $a_9^9[a_9^9]2<4\ [a_9]2<4\ a_9e$,

$\text{Re}\,\lambda_9\left[\lambda_9\right]_9 = \text{Re}\lambda_9\left[\lambda_9\right]_9 = -\frac{9}{9}a_9[a2\ ]_9 < 0.$

Let $\Delta=\max_{9<u<\infty}\left|h_*\,[h_*]_9\beta_9(a_9) - \mu_9(a_9)\right|$. Then$|n_9(t)|$ ⍰⊿∀t ⍰0. Hence we obtain the following result.

**THEOREM 4**

Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_9^9 <4\ a_9\,e$ and $[a_9] > 1\Delta$,or $[a_9^9]$ ⍰4$[a_9]e$ and $(a_9 - \sqrt{(a_9^9)_9 - 4(a_9)_9 e) > 1\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+9}\right)$ and $\beta_9(a_9) = ae^{-ma}$ Then $\Delta=[\mu_0] + 2$ Simple calculation leads to

$$\int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9) + d\tau_9 da} = \frac{1}{(m + \mu_0 + 1)^9} + \frac{1}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_9 + [\mu_0]$ 0.5. Assume$[f_9](x_9)= K_9 x_9$ Then the hypothesis (H9) is satisfied We take the conversion rate function $g_9(x_9)= [0x_9]/(2 + x_9)$; the functional responses $[[f_9](x) = [0]\, x_9/(1 + x_9)$and $h_9(x_9)=(1+\omega_9 x_9)/(1 + x_9)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_9 + [\mu_0] + [K][P_0])}{[m_9]_9 + (\mu_0)_9 + [K][P_0])_9^3\ [K]_9[P_0]_9^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_9$i e $[P_0] < b/(\delta_9 - b_9)$then the equilibrium$(P_0)$ $[P_0]_9(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_9$

$[f_{9*}] = K_9\ P^*\ [f_{9*}] = \frac{\sigma Q_9^*}{9+Q_9^*}$ $[h_*] = \frac{9+\omega_9 Q_9^*}{9+Q_9^*}$

$[f_{9*}] = K_9\ [f_{9*}^9] = \frac{\sigma_9}{(9+Q_9^*)^9}\ [g_*^9] = \frac{\sigma_9}{(9+P_9^*)^9}$ $[h_*^9] = \frac{\omega_9 - 9}{(9+Q_9^*)^9}$

In order to have (H3) and (H4) hold,$m, \mu_0, K, b, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_9}{(m+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9+b_9}+\sigma)^9}+\frac{1\omega_9}{([m_9]_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9+b_9}+\sigma_9)^3}<1$$

And

$$\frac{1}{(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9})^9}+\frac{1\omega_9}{(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9}+(m_9+[\mu_0]_9+\frac{K_9 b_9}{\delta_9-b_9}+\sigma_9)\,Q_9^*)^3$$

$$=m_9+[\mu_0]_9+3+\frac{K_9 b_9}{\delta_9-b_9}(1+Q^*)(1+\omega_9 Q_9^*$$

Now

$$\begin{bmatrix}1\\1\end{bmatrix}-4\begin{bmatrix}a_9\end{bmatrix}e=\left(\frac{1K_9 b_9}{\delta_9-b_9}\right)^9+\frac{4K_9 b_9\sigma_9 Q_9^*}{(\delta_9-b_9)(9+Q_9^*)}+\frac{\sigma_9^9 Q_9^{*9}}{(1+Q_9^*)^9}$$

$$\frac{4\,b_9\sigma_9(\delta_9-b)Q^*}{\delta_9(1+Q_9^*)^9}-\frac{4\,(\delta_9-b_9)^9(1-\omega_9)Q_9^*}{\delta_9(1+\omega_9 Q_9^*)(1+Q_9^*)}P^*(0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_9 \, ae^{mu}_9 \quad [\mu_0]_9 - \frac{a_9}{a_9+9}$$

Hence if $[a^9_9 > 4(a_9)_9 e$ and

$$[a_9] = \frac{1K_9 b_9}{\delta_9 - b_9} + \frac{\sigma_9 Q^*_9}{1 + Q^*_9} > 1\Delta$$

Some of the examples to which the model has been applied are:

**New look at the galaxy-shaping power of black holes**

(477) "It creates (eb) these bubbles, and the bubbles keep the plasma hot.

(478) That's what prevents galaxies from (e) becoming even bigger than they are now."

(479) Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.

(480) In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.

(481) Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.

(482) "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.

(483) "It's like nature's thermostat, which keeps these galaxies from (e) growing.

(484) If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."

**A new look at the galaxy-shaping power of black holes**

(485) An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the

(486) McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.

(487) Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.

**Doomed satellite missions**

(488) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.

(489) At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.

(490) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.

(491) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.

(492) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.

(493) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people

(494) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."

(495) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.

(496) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(497) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.

(498) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.

(499) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.

(500) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.

| |
|---|
| (501) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well. |
| (502) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma. |
| (503) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6. |
| (504) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus. |
| (505) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (506) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (507) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (508) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (509) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (510) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (511) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (512) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (513) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (514) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (515) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (516) We know dark matter exists only from (e) its gravitational effects. |
| (517) It holds together galaxies and (e&eb) galaxy clusters. |
| (518) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (519) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |

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| (520) | They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (521) | X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (522) | They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (523) | This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (524) | An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |
| (525) | Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (526) | "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (527) | The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (528) | Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (529) | Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (530) | "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (531) | Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (532) | "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (533) | Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) |

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| consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories. |
| (534) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen. |
| (535) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer. |
| (536) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants. |
| (537) These were treated as (=) "aware" and "unaware" trials, respectively. |
| (538) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it. |
| (539) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another. |
| (540) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion. |
| (541) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness. |
| (542) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said. |
| (543) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness." |
| (544) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner. |
| (545) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain. |
| (546) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said. |
| (547) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

(548) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter.

(549) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined.

(550) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions.

(551) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.

(552) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.

(553) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.

(554) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.

(555) Noda and his colleagues instead studied the change in emissions over (e&eb) time.

(556) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.

(557) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).

(558) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.

(559) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.

(560) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(561) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking "palm tree" organization.

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| (562) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society |
| (563) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible? |
| (564) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem. |
| (565) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other. |
| (566) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible. |
| (567) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges. |
| (568) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century. |
| (569) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching. |
| (570) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom. |
| (571) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth. |

**Spheres on trees**

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| (572) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration. |
| (573) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement. |
| (574) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum. |
| (575) By visualizing the problem in this way, the researchers noticed that these particular graphs don’t (e) have lots of branches extending from lots of (e) other branches. |
| (576) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure. |

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| (577) The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network. |
| (578) As a result, it provides (eb) an important clue for finding the global minimum. |
| (579) It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum. |
| (580) This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large. |

**Thompson's problem 2**

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| (581) By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society |
| (582) The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima. |
| **(583)** The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node. |

**Implications for Thomson's and Smale's problems**

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| (584) Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before. |
| (585) "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around." |
| (586) As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem. |
| (587) "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view." |
| (588) In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges. |
| (589) From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear. |
| (590) "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK. |
| (591) "For larger numbers, there will likely be (=) multiple funnels. |
| (592) Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum." |

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| (593) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem. |
| (594) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said. |
| (595) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org |

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| The nontrivial equilibrium is stable. |
| IF we define $[a_9]_9 = [f_{9*}]_9 + [f_{9*}]_9 + [f_{9*}]\,P^*_9$ |
| $$[a_9]_9 = \left[f^{9}_{9*}\right]P^*_9 - \frac{h^9_*}{h_*}P^*(0)$$ |
| $$e=[g^9_*]_9\,Q^*_9$$ |
| $$\eta_9(t) = \left[h_*\right]_9\beta_9\left(\theta_9(t)\right) - \mu_9\left(\theta(t)\right),$$ |
| The system(32)can be written as the following system: |
| $$[M(t)]\left[a_9\right]M_9 - [a_9]_9W_9 - [n(t)]M_9$$ |
| $$w_9(t) = eM_9$$ |
| Write |
| $$A=\begin{pmatrix}-[a_9]_9 - [a_9]_9 \\ e \qquad 0\end{pmatrix}$$ |

$$B=(t)=\begin{pmatrix} [\eta(t)]_9 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_9 \\ W_9 \end{pmatrix}$$

The solution of (32)is

$$X_9(t) = e^{A,t}[X]_9(0) + \int_0 e^{At}$$

**Note: There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and**

**sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by**

**term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation**

**alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation,**

**appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module**

**LEN NUNNEYS**

## Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

## Module numbered one

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems

and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1965); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be referred to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 6985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_6}{dt} = [N(t)]_6[(F_6)(R(t)]_6 - [M]_6[R(t)]_6$$

$$\frac{d[R(t)]_6}{dt} = [B]_6\,[(R)_6(t-T_6)] - [D]_6[(R(t)]_6 - [N(t)]_6\left[G\left(R(t)\right)_6\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_6}{dt} = a_6[N(t)]_6 - a_6[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_6 = a_3[N(t)]_3 - a_4[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_6 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

irst, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.6b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The

regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 6a is at the

**TIME DELAY**

FIG. 6. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.6$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**DELAY DUE TO PREDATOR MATURATION**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_6}{dt} = N_6(t - T_6)F_6(R_6(t - T_6) - N_6(t)M_6((R_6(t) \qquad (5)$$

$$\frac{dR_6(t)}{dt} = B_6\big(R_6(t) - D_6(R_6(t)\big) - N_6(t)G_6(R_6(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_6}{dR_6} + \frac{N_6 dG_6}{dR_6} - \frac{dB_6}{dR_6} > \frac{N_6 G_6}{2M_6}\left[\frac{dF_6}{dR_6} + \frac{dM_6}{dR_6}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_6/dR_6$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_6$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_6$) and the other derived from the non-delayed or "immediate" processes($I_6$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

ABSOLUTE STABILITY

$$L_6 = \frac{N_6 G_6}{2M_6}\frac{d(6nF_6)}{dR_6} (\geq 0) \qquad (6a)$$

$$I_6 = -\frac{N_6 G_6}{2M_6}\frac{d(6nM_6}{dR_6} (\geq 0) \qquad (6b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_6 > L_6 - I_6$, must be modified ahe $dM_6/dR_6$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_6 > L_6 - I_6, \quad \text{when} \quad P_6^2/4 > I \qquad (8a)$$

$$P_6 > L_6 - I_6 + \frac{(I_6 - P_6^2}{(L_6 - I)} \quad \text{when} \quad P_6^2/4 < I. \quad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

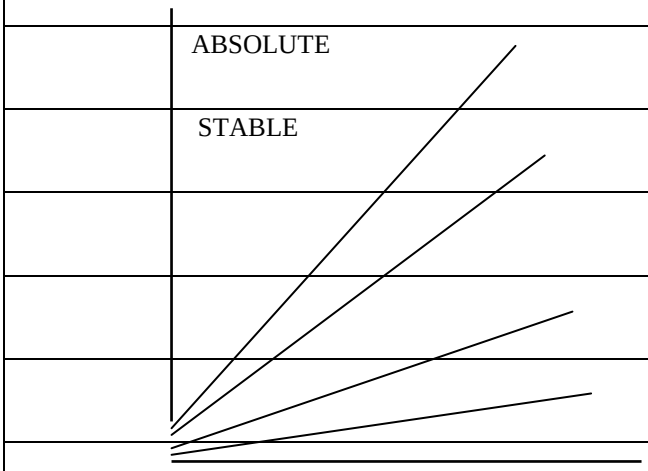


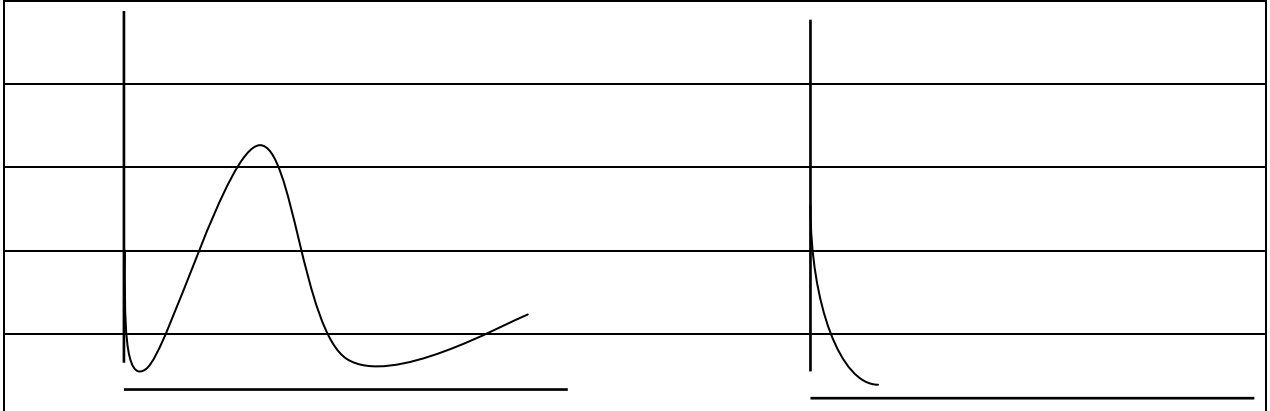

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

Can never completely offset the lagged effect,$L_6$; condition (8b) ensures that the total debit regulation, $P_6$, must always be positive even when $I_6$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_6$ and $L_6$. As $I_6$ increases relative to $L_6$ the zone of stability is enlarged; however,even when $I_6$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_6$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



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Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (6). Parameters in each case M=6; $dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6 = 0.05$ For (a) $NGdF/dR = 0.3, -N_6G_6dM_6/dR_6 = 0.6$; for (b)$NG_6dM_6/DR_6 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_6t}{dt} = N_6(t - T_6)F_6\big(R_6(t - T_6)\big) - N_6(t)M_6\big(R_6(t)\big)$$

(9)

$$\frac{dR_6t}{dt} = B_6\big(R_6(t - T_6)\big) - D_6\big(R_6(t)\big) - N_6(t)G_6\big(R_6(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) and the effective prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we

could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

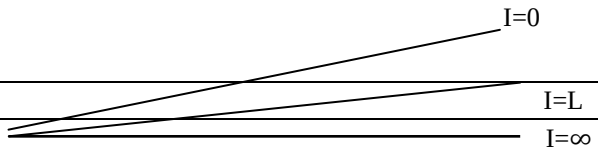
Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_6/dR_6$ is a measure of the control acting through the debit delay and setting $dB_6/dR_6 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_6/dR_6$ is a measure of the control acting through the credit delay although setting $dF_6/dR_6 = 0$ does not negate the effect of a predator delay; the value of $M_6$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_6/dR_6$and normalized by $M_6$ (see6a). The absolute stability of such a system is summarized in Fig. 6 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

TOTAL

ABSOLUTELY STABLE

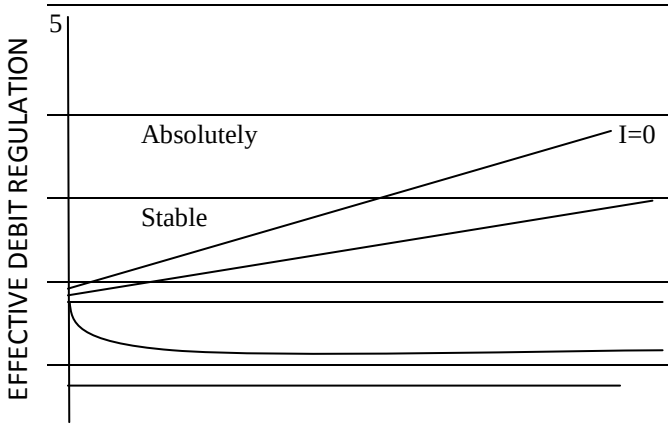



DELAYED DEBIT EFFECT

FIG6. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_6/dR_6 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.



-50 10

DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_6/dR_6 + N_6 dG_6/dR_6 + dB_6/dR_6)/M_6$). The immediate ($I_6$) and lagged ($L_6$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_6$(see 6b), is increased relative to the lagged part, $L_6$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_6 > I_6$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_6$is somewhat large than $I_6$ then the condition constraining prey regulation of recruitment ($dB_6/dR_6 < 0$) is

$$\frac{dD_6}{dR_6} + N_6\frac{dG_6}{dR_6} + \frac{dB_6}{dR_6} > \frac{(L_6 - I_6)M_6}{2} \qquad (10)$$

If $L_6 < I_6$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_6$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_6$

The - $D_6$ positions are defined by letting $\lambda_6 = iZ_6$ and separating (A1.2) into real and imaginary parts:

$$Z_6^2 + V_6 Z_6 \int_0^{\infty} \sin(Z_6 T_6) f(T_6) dT_6 - (X_6 + Y_6) = 0$$

$$Z_6 \left( U_6 - V_6 \int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6 \right) = 0$$

Noting that $X_6 + Y_6 = 0$ is we assume X+Y=0.It follows that $Z_6 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_6 = V_6 \int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6$$

If $U_6 > |V_6|$then given the constraints of $f(T_6)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A6.5)remains

greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

## ABSOLUTE STABILITY WITH A DELAY DUE TO CREDIT CONSOLIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_6^2 + (U_6 - V_6 + M_6)\lambda_6 - M_6\lambda_6 e^{-T_6\lambda_6} + [(U_6 - V_6)M_6 + Y_6] - [((U_6 - V_6 M_6 -]e^{-T_6\lambda_6} = 0$$

(For definitions see Appendix6) We follow the technique outlined in Appendix 6 and let $\lambda_6 = iZ_6$ and separate (A2.6) into real and imaginary parts

$$Z_6^2 + M_6 Z_6 \sin(Z_6 T_6) - (A_6 M_6 + Y_6) + (A_6 M_6 + X_6)\cos(Z_6 T_6) = 0$$

$$(A_6 M_6) Z_6 - M_6 Z_6 \cos(Z_6 T_6) + (A_6 M_6 - X_6)\sin(Z_6 T_6) = 0$$

Where $A_6 = U_6 - V_6$Using (A2.3), substitute for $Z_6^2$ and $Z_6$in (A2.2):

$$W_6^2(X_6^2 - 3M_6 A_6 X_6 - M_6^2(X_6 - Y_6) + 2M_6^3 A_6 + 2A_6^2 M_6^2]$$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_6 = 1 - \cos(Z_6 T_6)$.

In the absence of a delay the system is stable for $A_6 > 0,$ and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_6 > (X_6 - Y_6)/2\ M_6$ Thus a necessary condition for absolute stability is

$$A_6 > [\ X_6 - Y_6/2\ M_6$$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_6$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_6$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_6, A_6) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_6 \le 2$ *is* of interest First consider the extreme values

$$Q_6(0, A_6) = A_6^2\ (X_6 - Y_6)$$

$$Q_6(2, A_6) = (A_6 + 2\,M_6)^2\,(2A_6 M_6 - (X_6 - Y_6))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_6}{\partial A_6} = W_6^2\left(2M_6^3 + 4A_6 M_6^2 - 3M_6 X_6\right) + W_6\left(3M_6 A_6^2 - 2A_6 X_6 + 2M_6\left(2X_6 + Y_6\right)\right)$$

$$+2A_6\left(X_6+Y_6\right)$$

$=2M_6^2W_6^2(M_6+\mathrm{A}_6)+M_6A_6^2W_6+\left(2A_6M_6-X_6+Y_6\right)(A_6W_6+M_6W_6^2$

$+(2-W_6)((2M_6X_6+M_6Y_6)W_6+A_6(X_6+Y_6)),$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_6$quadratic in $\mathrm{W}_6$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_6$has no minimum in the range 0≤ $W_6$ ≤2.Defining

$Q_6\left(W_6,A_6\right)=a_6W_6^2+c_6 where\ a_6\ b_6\ and\ c_6$ Are functions of $A_6$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_6\geq b_6\geq 0$$

And

$$b_6^2\geq 4a_6c_6$$

$$Case(i):Y_6=0$$

$Let\ Y_6 = 0$ And try the substitution A= $X_6 = /2M_6\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_6 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_6 X_6 /$ $2M_6$) is positive for *W<2. Since* $Q_6$ increases with $A_6$(see A2.9), increasing $A_6$will generate no further solutions of $Q_6$=0. Hence the condition A>X/2M defines absolute stability when $Y_6 = 0$.

$Case\ (ii)$: $Y_6 = y_6\ X_6$ ,

$Let\ Y_6 = y_6\ X_6$And try the substitution $A_6 = X_6/(2 + A_6\ y_6)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_6/(2 + n_6 y_6)$ we have

$$W_6^2\ y_6\left[X_6 n_6 \frac{(1 + n_6 y_6)}{(2 + n_6 y_6)} + M_6^2\left(2 - n_6 + n_6 y_6\right)\right]$$

$$-W_6\ X_6\left[\frac{X_6(1 + n_6 y_6)}{M_6^2(2 + n_6 y_6)^2} + 2\ y_6\left(n_6 - 1\right) + \frac{X_6^2(6 + y_6)}{M_6^2(2 + n_6 y_6)} = 0.\ (A2.12)\right.$$

Since $n_6$ and $A_6$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_6$

Defining $E_6 = (2 + n_6 y_6), the\ condition (A2.66?) can\ be\ written\ as$

$[\frac{X_6(E_6 - 6)}{E_6^2}]^2[\frac{X_6^2}{M_6^2} - 4\ y_6\ E_6^2\frac{X_6}{M_6^2} + 4y_6(n_6 - 2)E_6^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in ($X_6 / W_6^2$) can be expressed as

$2E_6[y_6E_6+(y_6E_6(6+y_6)(E_6-n_6))^{6/2}]>\frac{X_6}{M_6^2}>2E_6[y_6E_6-y_6E_6(6+y_6)(E_6-n_6))^{6/2}]$

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Hence if $X_6/M_6^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_6$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_6-6)}{E_6^2}\left(\frac{X_6}{M_6^2}\right)^2-2y_6\frac{(E_6+n_6^2y_6)}{E_6}\left(\frac{X_6}{M_6^2}\right)-4y_6(E_6-n_6)$$ 0

And has the solution

$\frac{X_6}{M_6^2}>\frac{E_6}{(E_6-6)}[y_6(E_6+n_6^2y_6)+[y_6^2(E_6+n_6^2y_6)^2+4y_6(E_6-6)(E_6-n_6)]^{6/2}]$

Given $E_6>n_6$ this lower limit to $X_6/M_6^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_6/M_6^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_6$ in the appropriate range and this limit is further lowered by (A2.64)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_6<n_6$ $(ie.,2A_6M_6^2<X_6-Y_6)$ then (A2.6)and

(A2.8)show that $Q_6$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_6}{M_6^2} > 2E_6[y_6E_6 - (y_6E_6(6 + y_6)(E_6 - n_6)]^{6/2}]$$

This equations throws light on the individual general ledger and Nature's general ledger as we call it : see various papers on the issue. powers, which are the energies of Matrka we have already discussed, are not the only obstacles the yogi must overcome. He must, for example, also resist the temptation to rest content with the miraculous yogic powers (siddhi) he acquires in the course of his spiritual develop- ment. Again to do this he must practice Yoga. Similarly, in order to pervade the Fourth State gradually through the other states in the manner proper to practice at the Individual level, the method is the same. He must practise the higher yoga of the Tantras which, turning his mind inwards and freeing it from discursive representations, allows him to penetrate into the Supreme Principle. 312 Once the yogi has attained this contemplative state, his main problem is to make it permanent. In the introverted state the gross external movement of the breath is suspended and with it the activity of the intellect, mind, individualised consciousness, powers of the senses and the ego. 313 When the yogi rises out of this state, he is liable to fall again into the lower order of creation generated by Maya if he does not maintain his awareness of the higher reality he has experienced and allows his awakened, illumined insight to be obscured by the dream-like vision of thought-constructs. 314 Naturally, the yogi must rise out of the introverted condition of suspension. It is inherent in the very nature of reality that it should move out of itself. 315 Pure, universal consciousness initially transforms itself into the vital breath 316 charged with the impression iyasana) of the power of awareness attained through introversion. By attending to the pulse (spanda) of the breath as it moves out of the absolute, the

yogi can develop an intuitive sense of the inherent unity of all he will perceive in the mental and physical spheres created by the outpouring of consciousness. In this way he realises that his own nature is everywhere present in all he perceives and that all things thus reside within him. Blessed with this insight his consciousness remains free and unlimited even at the individual level where the breath, mind, senses and body are active. If the yogi fails to do this, he finds himself once again beset by the strictures of his embodied existence and must, as before, try to pervade all his other states of consciousness with the aesthetic delight (rasa) and wonder (camatkara) of the Fourth State he experienced in contemplation. Again this means that he must strengthen his pure, empowered awareness that his universal nature manifests as all things. 317 In this way he discovers Siva's presence in every sphere of individualised consciousness ranging from the breath to outer objectivity. The yogi's mind then becomes tranquil and undistracted because wherever it may wander, the yogi perceives only Siva, his authentic nature. 318 Consciousness is thus freed of all external referents and the yogi's subjectivity is purified of all identification with the body or anything else that belongs to the objective sphere. The yogi then becomes detached from the opposites of pleasure and pain and is transcendentally free (kevalin). 319 yogis are again, however, liable to fall if he allows himself to get entangled in the play of opposites. This fall is more serious than the others because, although he is caught by the confining restrictions of individual- ised consciousness as before, he is now also affected by karma. Fleeing from pain in the pursuit of pleasure he is bound to act (karma) to minimise one and maximise the other and so is thrown down to the lowest level of embodied subjectivity (sakala). In order to regain his lost state, he must ascend gradually, by Siva's grace, from one order of subjectivity to the next and so free himself progressively of the limitations of the lower levels to gain the greater freedom and expansion of the higher. As he progresses, the objective sphere also evolves from the grossest perceptions of physical objects outside the lowest order to subjectivity, through to the subtler inner, mental

perceptions to finally reach the order of subjectivity that contains objectivity within itself and is free to externalise it at will. 320 The degree to which this process develops depends, as before, on the yogi's awareness of the Fourth State. In consonance with the general principle that the remedy should suit the defect, the yogi is instructed to seek this higher state of consciousness in the wonder (camatkara) or delight (ananda) he feels in moments of intense physical pleasure. At first he experiences this subtle consciousness for an instant in the subjective sphere. If he manages to catch hold of it, it becomes more intense as the cognitive and objective spheres are also gradually pervaded and vitalised by it. Occasions for this practice are, for example, the sense of satisfaction one feels after a good meal or the aesthetic delight one experiences when listening to good music or the pleasure of sexual union with the Tantric consort or even solitary sexual excitation. In these moments of delight the yogi can penetrate momentarily into his own authentic Siva-nature (sambhavavesa) through the empowered contact (saktasparsa) m he makes with it in the freedom of the pure subjectivity of the Fourth State. 322 If the yogi develops his awareness of this higher level of consciousness and maintains it, he eventually experiences it constantly. 323 Clearly, what prevents the yogi from attending to his state of consciousness rather than the circumstances which induce it is the craving for pleasure (abhila§a) born of ignorance — the source of every impurity which clouds consciousness. Craving directs the yogi's attention towards outer, worldly things and so he is caught in the net of thought- constructs. 324 To free himself of his worldly desires and reverse this binding extroversion, the yogi must eradicate its cause. To be freed of all the ups and downs of the path and no longer be tormented by the possibility of a fall, the yogi must see reality perfectly and completely. This insight is itself liberation and the moment it dawns the yogi is instantly freed. This sudden realisation is the goal of Tantric Yoga. E Tantra declares: "He who perceives reality directly, even for the brief moment it takes to blink, is liberated that very instant and never reborn again." 325 Although the yogi's body and mind

continue to function as before, they are like mere outer coverings 326 which contain, but do not obscure, the mighty, universal consciousness which operates through them. The yogi's body is (=) the universe, the senses (=) the energies that (eb, eb+)vitalise it, his mind (=) Mantra, the rhythm of his breath (eb) the pulse of time and his inner nature pure, dynamic consciousness. Raised above all practice, and hence all possibility of falling to lower levels, the yogi realises that he has always been free 327 and that his journey through (e&eb) the **dark land of Maya was nothing but a dream, a construct of his own imagination** The channels of the senses by which the yogi enjoys the pleasure of the objects of sense are those through which he fills the three worlds with the consciousness emitted from his own Heart. 78 **The Doctrine of Vibration T**he senses are outer, subsidiary cycles of conscious energy (anucakra), revolving around the core or primary cycle (mukhyacakra) of pure consciousness. They scintillate with the diversity of sensations and resonate in the field of their awareness like the feathers of a peacock unfurled and folded back in the ecstasy of its dance. 79 Ultimately it is Siva Himself Who, for His own pleasure, 80 extends and withdraws them, 81 not the individual soul. The outward movement of the senses is the creative flow of consciousness from the subject who, like an ocean, is the source of the waves of their activity and the final resting place of the flow of the rivers of sensations and thought-forms that unfold through the arising of the power of each of the senses. 82 The Stanzas on Vibration declare: Indeed the individual soul does not activate the impulse of the will [which directs the body's activity] by himself alone but through his contact with [his] own [inner] strength {bald) made in such a way that he identifies with it [thus acquiring its power]. 83 The close contact between the senses and universal consciousness vitalises them and sets them in motion. They share in the properties of consciousness, just as a heated iron ball can burn, cook or give out light. 84 Spanda doctrines does not agree with the Sarnkhya view that the individual soul (purusa) alone stimulates activity in the objective sphere (prakrti) by the contact he makes with it. The soul can do

this only because universal consciousness animates him. Moreover, consciousness, which has the power to impel this activity, cannot itself be inactive as the Sarnkhya maintains. The influence it has on lifeless matter implies an active exertion. 85 This activity is the subtle, inner outpouring of the Self (svatmocchalatta) within the Self. Outside the confines of time and space, in the centre of the circle of the senses, this inner activity initiates the expansion (prasara) of the flow of conscious energy. As the pure exertion (udyoga, udyama) of Being (bhava), it is the source of the initial unfolding (unmesa) of Becoming. 86 Ksemaraja explains: [By] 'exertion' (udyama) is [meant] the emergence of the supreme [level of] intuition (pratibha) which is the sudden outpouring of the [immanent and transcendental] unfolding of the reflective awareness of pure consciousness. 87 The direct, intuition (pratibha) of reality, free of thought-constructs, is the supreme power (parasakti) of consciousness. The yogi can gain Divine **Body and Sacred Circle of the Senses 153** this pure intuitive awareness by grasping the initial expansion (unmesa) of consciousness which engenders the emergence of thought and the activity of the senses. It can be discovered in the subtle Centre, between one thought and the next. The Stanzas on Vibration teach: The expansion of consciousness that takes place when one is engaged in a single thought should be known to be the source from whence another arises. One should experience that for oneself! 88 This does not mean that the unfolding expansion of consciousness {unmesa) is an intermittent, transitory occurrence. 89 As the source of all thought, including the notions of past, present and future, 90 it stands outside the confines of time. Consciousness is in a state of perpetual expansion. It is an ever developing process, a tiny part of which we observe as the flow of events around us and the thoughts and feelings within us. The totality of this process can never be grasped objectively. Objectivity is a state of fragmentation (bheda) in which discrete elements are discernible and mutually distinguished as individual, specifically definable entities. It is the realm within which thought and language operate. To grasp reality in its completeness, we must go

beyond the partial representations of thought and speech. To experience the primordial source and basis of all things, we must pierce through the outer periphery of thought and plunge into the Centre, to discover the instant in which thought, and with it, the sense of diversity, initially emerges. Ksemaraja cites a practice from the Vijnanabhairava through which the yogi can achieve this state of awareness: Checking [the movement of] attention (cit) once it has quit its object, it must not [be allowed] to move to any other; contemplation (bhavana) then blossoms forth by [experiencing] the state in the centre between them (tanmadhyabhava). 91 Left unattended and undisciplined, awareness spontaneously shifts from one object of thought to (e&eb) the next. The transition from one to the other entails a movement through (e&eb) a state of pure indeterminate (nirvikalpa) awareness divested of all thought-forms and perceptions. The yogi who succeeds in checking the movement of his attention, experiences within himself the subsidence of thought into the emptiness of the Centre which unfolds full of the powerful pulsation of consciousness within which all objectivity merges and becomes one 92 in the trans- cendental outpouring (lokottarollasa) of its aesthetic rapture. 93 turning in on himself, the yogi must exert his power of awareness to (e) discern the vitality of (e) his own blissful Spanda nature. Just as a distant object becomes (=) more clearly visible by (e) exerting an effort to (e) perceive it, in the same way, through (e&eb) practice, the yogi's own Spanda nature becomes (=) progressively more evident to him, 94 The effort the yogi exerts to free himself of ignorance is essentially one with the force of awareness that impels the senses and mind. This inner exertion is the basis of every means to liberation. Ultimately bestowing the bliss of self- realisation, it carries the yogi along the path to salvation with all its twists and turns like wind which blows a cloth in various directions. 95 As the basis and ultimate goal of all yogic practice, this exertion is the pure wonder (camatkdra) inspired by the realisation of the essential identity of consciousness (cit) and being (sat). 96 As such it is the divine radiance (sphuratta) of the Heart of one's own consciousness 97 where the powers of

will, knowledge and action merge in the harmony (samarasya) of bliss. 98 Thus the highest form of worship (varivasya) is the contemplation of the inner strength (bald) or universal exertion of one's own true nature. 99 "One should examine [that reality]," says Ksemaraja, "with the exertion which is Bhairava's nature and is true service (sevana) to the integral inner nature [of all things] and the unfolding of the activity of one's own vitality (ojas) which withdraws all duality [back into consciousness]." 100 Mindful of the eternal joy (ahlada) of his inner strength, the yogi must give up all desire for the fleeting moments of petty pleasure (ksanikasukhalava) ]0] he may glean from outer objects. He must even give up the desire to achieve enlightenment and a reality (tattva) which can never be realised through a personal act of will. 102 In short, he must die to himself in an act of profound faith (sraddha) and adoration (adara) to find rest in the supreme reality free from the binding dichotomy of means and objectives (upeyopayabhava) — the reality of his perfected effort. The yogi, now full of devotion for this reality unfolding within himself and the universe (antarmukhatattva), merges into it and, by achieving the pure consciousness and bliss (cidananda) of the power (bala) of his own nature, achieves oneness. 103 The yogi is firmly established in his true nature by recognising the creative unfolding (unmesa) of consciousness through which he is able to know all things, whether in the present, past or future. 104 Master of the vitality which gives life to his mind and body, he can strengthenit directly without need of food or drink, 105 or quit it at will to vivify Divine Body and Sacred Circle of the Senses 155 that of another. 106 Experiencing the eternal satisfaction of his own uncreated nature, 107 the yogi is freed of disease and the infirmities of old age. 108 Inspired by this vitality, he has the inner strength to overcome the severest hardships and obstacles in the way of his ultimate goal. 109 The Stanzas on Vibration explain: Lassitude ravages the body and this arises from ignorance but if this is eliminated by an expansion of consciousness (unmesa), how can [ignorance], deprived of its cause, continue to exist? 110 Doubt (sanka) is the source of every spiritual ailment. 111 A

man in doubt about his true nature, and the path to follow by which he may come to realise it, is constantly overcome by difficulties. 112 "Due to ignorance" teaches the Sarvacaratantra, "the deluded man is in doubt and thus [suffers the cycle] of birth and death." 113 Caught between conflicting alternatives, the mind shifts from one to the other becoming more entangled in its own thoughts as it does so. Doubt thus contracts consciousness 114 sullying it with the turbulence of confused thinking. 115 At once both the root of "the ancient tree of transmigration" and the first sprout of its seed, 116 doubts deprives man of the innate bliss (sahajananda) of his own nature. 117 Like a thief it steals away the wealth of true knowledge and reduces man to an imagined state of poverty in which he feels hemmed in on all sides by constraints and limitations. 118 Ignorance and its sister, doubt, are the essence of all impurities (mala) which sully consciousness. To the degree in which the yogi is cured of this ailment, namely, the lassitude (glani) of doubt, his true nature becomes manifest, just as raw gold when heated is freed from dross. When the yogi recognises the all-powerful expansion (unmesa) of his consciousness, he is projected beyond the realm of relative distinctions between virtue and vice. Thus overcoming all doubt as to what he should or should not do, he penetrates into the pulsing Heart of Bhairava's consciousness. 119 As the Stanzas on Vibration say: An individual who [though] desirous of doing various things [but] is incapable of doing them due to his innate impurity, [experiences] the Supreme state when the disruption (ksobha) [of his ego] ceases. Then [the soul realises] that his [true] uncreated nature is [universal] agency and perceiving subjectivity and so he knows and does whatever [he] desires. 120 the individual soul is(=) a combination of (e) conscious (cit) and 156 The Doctrine of Vibration unconscious (acii) elements 121 brought together when he allows the power of his authentic nature to be obscured by (e) his own impurity. This impurity is (=) three-fold. The first is technically called (=) the 'Impurity of Individuality' (anavamald). Due to (e) this impurity, the individual soul fails to (e) recognise his all-embracing fullness (purnata) and is disturbed by (e) his craving

for (e) experience (bhoga) based on (e) the unreasoned assumption that he is (=) incomplete (apurnd) and in need of (e) something outside himself. 122 Thus it manifests as (=) desire or attachment (rdga). ]23 At the same time it is (=) primordial ignorance, 'the cause of the sprout of (e) transmigration'. 124 As such, it works in two ways. Firstly, it deprives (e) individual consciousness of freedom. The yogi affected by (e&eb) this impurity is incapable of (e) entering or emerging out of the states of contemplation as he wishes. Secondly, this impurity renders (eb) the individual unconscious of his freedom, which for(e) the yogi means that he is (=) liable to lose consciousness of (e) his true nature when he is (=) not in a state of (e) contemplative absorption. 125 The second impurity is (=) called the 'Impurity of Maya' (maylyamald). This comes into operation when consciousness has been contracted by (e) the Impurity of Individuality. It limits (e) the power of knowledge operating in (e&eb) the individual soul through (e&eb) the inner and outer organs of (e) perception. 126 Due to this impurity he perceives (e) diversity everywhere. 127 Finally, the third impurity, that of Karma, comes into play when (e) the individual, deprived of (e) the freedom and knowledge of (e) unconditioned consciousness, acts (e&eb) in his limited way prompted by (e) desires and fears for (e) his personal gain. Disturbed by (e) these three impurities the soul cannot find rest in (eb) himself.

And

$$E_6 > n_6$$

Replacing $E_6$ and $n_6$ in these conditions leads to the form

If $A_6^2 > 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6$.

If $A_6^2 < 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6 .+ (Y_6 - A_6^2/2)^2/2M_6(X_6 - Y_6)$.

Note that by normalizing $A_6$ by $M_6$ and $X_6$ and $Y_6$ by $W_6^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_6^2 + (M_6 + U_6)\lambda_6 - (M_6 + V_6)\lambda_6 e^{2T_6\lambda_6} + (M_6 U_6 + Y_6)$$

$$-[M_6(U_6 + V_6) - X_6]\, e^{2T_6\lambda_6} + M_6 V_6\, e^{2T_6\lambda_6} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_6 = \mathrm{i}\, Z_6$ *and* separate

(A3.1)into real and imaginary parts

$$Z_6^2 + (M_6 + V_6) Z_6 \sin(Z_6 T_6) - (U_6 M_6 + Y_6) + (M_6(U_6 + V_6) - X_6)$$

$$\mathrm{X}\cos(Z_6 T_6) - 2M_6 V_6 \cos^2(Z_6 T_6) + M_6 V_6 = 0$$

$$(M_6 + U_6) Z_6 - (M_6 + V_6) Z_6 \cos(Z_6 T_6) + (M_6(U_6 + V_6) - X)$$

$$\mathrm{X}\sin((Z_6 T_6) - M_6 V_6 \sin(Z_6 T_6) = 0$$

Using (A.3.3) to eliminate $Z_6$ and $Z_6^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_6 =$

$1 - \cos Z_6 T_6$ and $A_6 = U_6 - V_6$ we obtain

$$4 M_6^2 V_6^2 W_6^4 + 2 M_6 V_6 W_6^3[(M_6 - V_6)^2 + A_6(M_6 + V_6) + 2(A_6 M_6 - X_6)]$$

$$+ W_6^2 [M_6 - V_6)^2 (2 A_6 M_6 + Y_6 - X_6) + (2A_6 M_6 - X_6)^2 - A_6 (M_6 - V_6)(2A_6 M_6 - X_6)$$

$$+ M_6 V_6 (A_6^2 + 4(X_6 + Y_6))]$$

$$+ W_6 [A_6^2 + 2X_6)(2 A_6 M_6 - X_6) + 2A_6 Y_6 (M_6 + V_6) - M_6 A_6^3] + A_6^2 (X_6 M_6 - Y_6) = 0$$

Being a quartic in $W_6$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_6 > 0$$

$$and$$

$$C_6 > (X_6 - Y_6)/2 M_6$$

$recalling\ that\ A_6 = U_6$ and defining $C_6 = U_6 + V_6$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_6 (0, U_6 V_6 = A_6^2 (X_6 + Y_6)$$

$$Q_6 (2, U_6 V_6) = (C_6 + 2M_6^2)^2 (2C_6 M_6 + Y_6 - X_6)$$

So that given (A3.6) the extreme $W_6$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_6$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_6$ the existence of a D-position at this point is

determined by letting cos $(Z_6 T_6) = -1\ in\ (A3.3)$

ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 e^{-\tau_6\lambda_6} + X_6 + Y_6 = 0 \qquad (A1.1)$$

Where $V_6 = dB_6/dR_6$ $U_6 = dD_6/dR_6 + N_6 dG_6/dR_6$ $X_6 = N_6G_6dF_6/dR_6$ and $Y_6 = -N_6G_6dM_6/dR_6$ (evaluated at the equilibrium). Note that $U_6X_6$ and $Y_6$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_6$ as $(T_6)$ . (It follows that $f(T_6)$ 0 for all $T$, and $\int_0^\infty f(T_6)dt = 1.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 \int_0^\infty e^{-\tau_6\lambda_6} f(T_6) + X_6 + Y_6 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MODULE NUMBERED TWO**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed

the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).. Haimovici's and Nunny's models have been applied to some of the Deleuzean and Heideggerian problems too : Scoring on that count are :

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time

delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t - T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if

only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

TIME DELAY

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

## MATURATION AND CONSUMMATION OF CONFEDERATION OF DEBIT AND CREDIT TYPES:

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t-T_7)F_7(R_7(t-T_7) - N_7(t)M_7((R_7(t) \qquad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7} + \frac{N_7 dG_7}{dR_7} - \frac{dB_7}{dR_7} > \frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7} + \frac{dM_7}{dR_7}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes $(L_7)$ and the other derived from the non-delayed or "immediate" processes$(I_7)$:as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_7 = \frac{N_7 G_7}{2M_7}\frac{d(7nF_7)}{dR_7} (\geq 0) \qquad (7a)$$

$$I_7 = -\frac{N_7 G_7}{2M_7}\frac{d(7nM_7}{dR_7} \ (\geq 0) \qquad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7 > L_7 - I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_7 > L_7 - I_7, \quad \text{when} \quad P_7^2/4 > I \qquad (8a)$$

$$P_7 > L_7 - I_7 + \frac{(I_7 - P_7^2}{(L_7 - I)} \quad \text{when} \quad P_7^2/4 < I. \qquad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

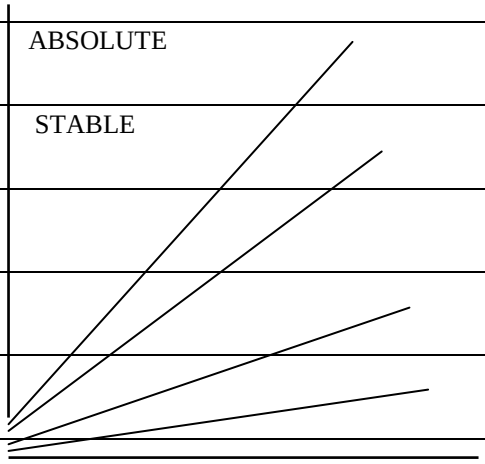



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit

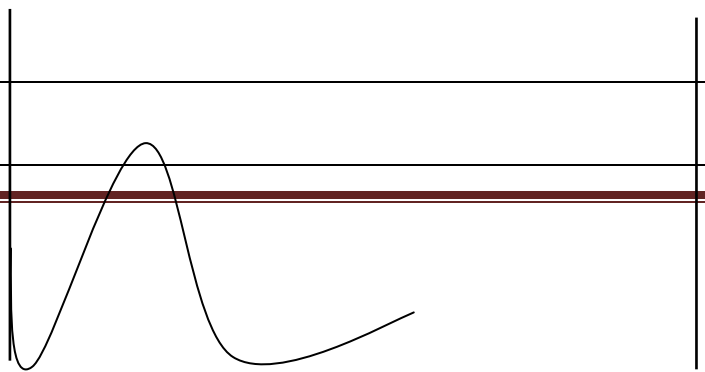
(like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

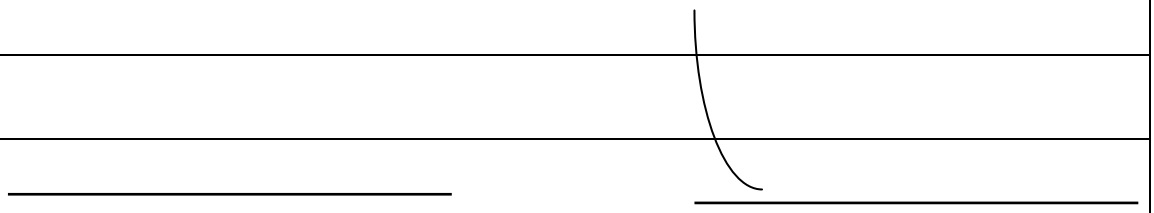
It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the





TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t-T_7)F_7\big(R_7(t-T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t-T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation $(dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7)$ and the effective prey regulation $(dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7)$ must be positive.

The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the

total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

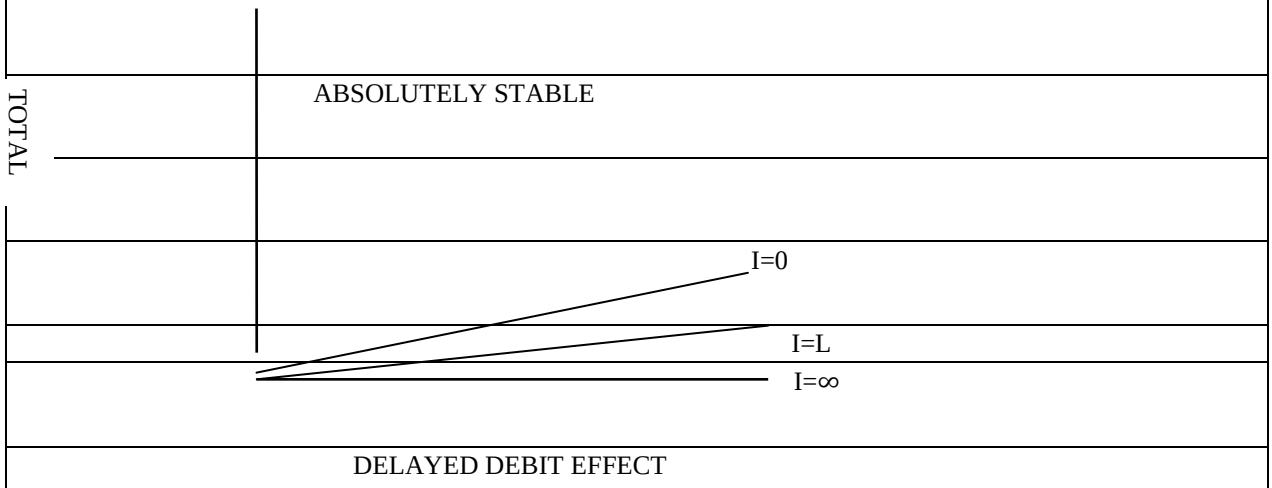
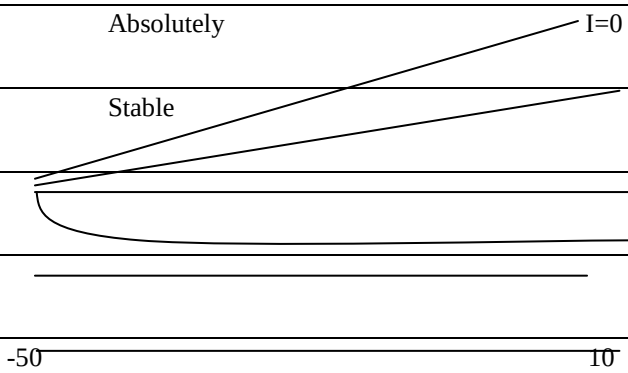



FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay.

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^\infty \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7\left(U_7 - V_7\int_0^\infty cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7\right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7\int_0^\infty cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$then given the constraints of $f(T_7)$, $(A1.5)$*can never be satisfied Co*nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7\lambda_7 e^{-T_7\lambda_7} + [(U_7 - V_7)M_7 + Y_7] - [((U_7 -$$

V 7M 7−]**e**−7 7λ 7=0

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$\boldsymbol{Z}_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7)\cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7)\sin(Z_7 T_7) = 0$$

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $\boldsymbol{Z}_7^2$ and $Z_7$in (A2.2):

$W_7^2(X_7^2$-3 $M_7\ A_7\ X_7$- $M_7^2(X_7$- $Y_7) + 2\ M_7^3\ A_7$+2 $A_7^2\ M_7^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_7$=1-cos $(Z_7 T_7)$.

In the absence of a delay the system is stable for $A_7 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7 > (X_7 - Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$$A_7 > [\ X_7 - Y_7/2\ M_7$$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_7, A_7) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_7 \le 2$ *is* of interest First consider the extreme values

$$Q_7(0, A_7) = A_7^2\ (X_7 - Y_7)$$

$$Q_7(2, A_7) = (A_7 + 2M_7)^2 (2A_7 M_7 - (X_7 - Y_7))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7 M_7^2 - 3M_7 X_7\right) + W_7\left(3M_7 A_7^2 - 2A_7 X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7+Y_7\right)$$

$=2M_7^2W_7^2(M_7+\mathrm{A}_7)+M_7A_7^2W_7+\left(2A_7M_7-X_7+Y_7\right)(A_7W_7+M_7W_7^2$

$+(2-W_7)((2M_7X_7+M_7Y_7)W_7+A_7(X_7+Y_7)),$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $W_7$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$has no minimum in the range $0 \le W_7 \le 2$.Defining

$Q_7\left(W_7,A_7\right)=a_7W_7^2+c_7 where\ a_7\ b_7\ and\ c_7$ Are functions of $A_7$we can define the condition for a minimum in the range 0-2

condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7 \ge b_7 \ge 0$$

And

$$b_7^2 \le 4a_7c_7$$

$$Case(i): Y_7=0$$

$Let\ Y_7=0$ And try the substitution A= $X_7=/2M_7\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}\,(W_7 X_7 / 2M_7)$ is positive for *W<2. Since* $\mathrm{Q}_7$ increases with $A_7$(see A2.9), increasing $A_7$will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

$Case\ (ii): Y_7 = y_7\ X_7$ ,

$Let\ Y_7 = y_7\ X_7$And try the substitution $A_7 = X_7/(2 + A_7\ y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + n_7 y_7)$ we have

$$W_7^2\ y_7\left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2\left(2 - n_7 + n_7 y_7\right)\right]$$

$$-W_7\ X_7\left[\frac{X_7(1 + n_7 y_7)}{M_7^2(2 + n_7 y_7)^2} + 2\,y_7\left(n_7 - 1\right) + \frac{X_7^2(7 + y_7)}{M_7^2(2 + n_7 y_7)}\right] = 0.\ (A2.12)$$

Since $n_7$ and $A_7$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7), the\ condition (A2.77?) can\ be\ written\ as$

$\left[\frac{X_7(E_7 - 7)}{E_7^2}\right]^2\left[\frac{X_7^2}{M_7^2} - 4\,y_7\ E_7^2\frac{X_7}{M_7^2} + 4y_7(n_7 - 2)E_7^3\right]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$2\,E_7[y_7 E_7 + (y_7\ E_7(7 + y_7)(E_7 - n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - y_7 E_7(7 + y_7)(E_7 - n_7))^{7/2}]$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0⁄$W_7$ ⁄2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7 - 7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2 - 2y_7\frac{(E_7 + n_7^2 y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right) - 4y_7(E_7 - n_7) \text{⁄0}$$

And has the solution

$$\frac{X_7}{M_7^2} > \frac{E_7}{(E_7 - 7)}[y_7(E_7 + n_7^2 y_7) + [y_7^2 (E_7 + n_7^2 y_7)^2 + 4y_7(E_7 - 7)(E_7 - n_7)]^{7/2}]$$

Given $E_7 > n_7$this lower limit to $X_7 / M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7 / M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$in the appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7 < n_7$ $(ie., 2A_7 M_7^2 < X_7 - Y_7$)then (A2.7)and (A2.8)show that $Q_7$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - (y_7 E_7(7 + y_7)(E_7 - n_7)]^{7/2}]$$

And

$$E_7 > n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2 > 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.

If $A_7^2 < 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7 + (Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$ and $X_7$ and $Y_7$ by $W_7^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSILIDATAN**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7\lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]e^{2T_7\lambda_7} + M_7 V_7 e^{2T_7\lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i}\, Z_7$ *and* separate

(A3.1)into real and imaginary parts

$$Z_7^2 + (M_7 + V_7)Z_7 \sin(Z_7 T_7) - (U_7 M_7 + Y_7) + (M_7(U_7 + V_7) - X_7)$$

$$\mathrm{X}\cos(Z_7 T_7) - 2M_7 V_7 \cos^2(Z_7 T_7) + M_7 V_7 = 0$$

$$(M_7 + U_7)Z_7 - (M_7 + V_7)Z_7 \cos(Z_7 T_7) + (M_7(U_7 + V_7) - X)$$

$$X \sin\left(\left(Z_7 T_7\right) - M_7 V_7 \sin\left(Z_7 T_7\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_7$ and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - cos\ Z_7\ T_7$ and $A_7 = U_7 - V_7$ we obtain

$$4 M_7^2 V_7^2 W_7^4 + 2 M_7 V_7\ W_7^3 [(M_7 - V_7)^2 + A_7 (M_7 + V_7) + 2(A_7 M_7 - X_7)]$$

$$+ W_7^2 [M_7 - V_7)^2 (2 A_7\ M_7 + Y_7 - X_7) + (2A_7\ M_7 - X_7)^2 - A_7 (M_7 - V_7)(2A_7\ M_7 - X_7)$$

$$+ M_7 V_7 (A_7^2 + 4(X_7 + Y_7))]$$

$$+ W_7 [A_7^2 + 2X_7)(2 A_7\ M_7 - X_7) + 2A_7 Y_7 (M_7 + V_7) - M_7 A_7^3] + A_7^2 (X_7\ M_7 - Y_7) = 0$$

Being a quartic in $W_7$ (A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7 > 0$$

$$and$$

$$C_7 > (X_7 - Y_7)/2\ M_7$$

$recalling\ that\ A_7 = U_7$ and defining $C_7 = U_7 + V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0, U_7 V_7 = A_7^2 (X_7 + Y_7)$$

$$Q_7\left(2, U_7 V_7\right) = \left(C_7 + 2M_7^2\right)^2 (2C_7 M_7 + Y_7 - X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1\ in\ (A3.3)$

## ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

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# Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price, University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss**

(1553) In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them.

(1554) Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors.

**Classical vs. relativistic black holes**

(1555) Something like a black hole exists within (eb) Newton's classical theory of gravity.

(1556) In that theory, an energy argument tells us that there is (=) an escape velocity vesc=2GM/R−−−−−−−√ from (e) the surface of any spherical object of (e) mass M and radius R

(1557) If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity.

(1558) Thus the condition for such an "unseeable" object is R<2GM/c2

(1559) In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape.

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| (1560) This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity. |
| (1561) To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it. |
| (1562) The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c . |
| (1563) This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed. |
| (1564) An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity. |
| (1565) Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable. |

**Stationary black holes**

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| (1566) In the Newtonian theory, gravity is described by (e) the potential $\Phi$. |
| (1567) Inside a spherical object the form of (e) $\Phi(r)$ depends on (=) the interior structure, but in the vacuum outside (e) matter the potential $-GM/r$ depends only on (=) the interior mass. |
| (1568) Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object. |
| (1569) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward. |
| (1570) This horizon radius $RH=2GM/c2$ is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory. |
| (1571) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution. |
| (1572) Radius RHhere actually means that the area of the event horizon is (=) $4\pi R2H$. |
| (1573) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution. |
| (1574) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity. |
| (1575) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own. |
| (1576) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole. |
| (1577) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density. |
| (1578) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed. |
| (1579) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole. |

(1580) In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter.

(1581) Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum.

(1582) In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior.

(1583) The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object.

(1584) Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum.

(1585) These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.]

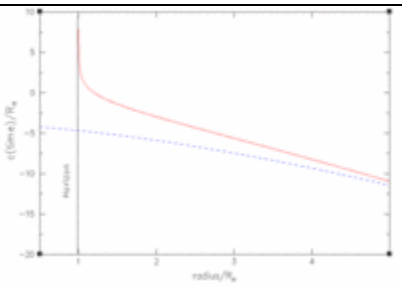
(1586) All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time.

(1587) But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?"

(1588) The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point.

(1589) One of these coordinates is (=) called the "coordinate time."

(1590) Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate.



(1591) Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole

(1592) At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole.

(1593) Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations.

(1594) For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree

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| (1595) An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events. |
| (1596) The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving. |
| (1597) Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle. |
| (1598) The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return. |
| (1599) But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon. |
| (1600) Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future. |
| (1601) The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?" |

**Black hole parameters**

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| (1602) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin). |
| (1603) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun. |
| (1604) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits. |
| (1605) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form. |
| (1606) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds. |
| (1607) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2. |
| (1608) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed. |
| (1609) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer. |
| (1610) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2) |
| (1611) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials. |
| (1612) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable |
| (1613) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite. |
| (1614) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole. |

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| (1615) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars. |
| (1616) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity. |
| (1617) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity. |
| (1618) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite. |
| (1619) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity. |
| (1620) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum. |
| (1621) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship. |
| (1622) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= (4πG2/c4)[(M+M2−a2−−−−−−−√ )2+a2]. |
| (1623) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored. |

**Dynamical black holes**

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| (1624) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape. |
| (1625) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, **a horizon can be born**. |
| (1626) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid. |
| (1627) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward. |
| (1628) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations. |
| (1629) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes. |
| (1630) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3. |
| (1631) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/ |
| (1632) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers. |
| (1633) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission. |

(1634) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.

(1635) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.

(1636) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.

(1637) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.

(1638) The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature $6\times10-8$ (Mo/M) K.

(1639) Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.

**Astrophysical black holes**

(1640) Black holes in our Universe can be grouped as: primordial black holes, stellar-mass black holes, and supermassive black holes. The first is highly speculative; the second and third are broadly accepted.

(1641) Primordial black holes of all masses are postulated to have formed from (e) quantum fluctuations in the early Universe, but those with (e&eb) mass less than around 1013kg would have already evaporated due to Hawking radiation.

(1642) No significant observational evidence has yet been found of (e) the existence of these objects [MacGibbon and Carr (1991)].

(1643) Stellar-mass black holes, ranging from a few to a few tens of solar masses, are (=) normal but rare endpoints in (eb) the evolution of massive stars.

(1644) When a star exhausts its nuclear fuel and cools, it must collapse unless (e) it is supported by nonthermal forces.

(1645) It is known that such nonthermal forces cannot (e) resist gravitational compression for (e) masses greater than around 1.5 to 3Mo. (The uncertainty is due to (e) uncertainty in our knowledge of nuclear physics at high densities.)

(1646) There are many stars far more massive than this, but in their death throes massive stars expel (eb) a great deal of their mass in supernova explosions

(1647) Even stars initially as massive as 20-30Mo may blow off enough mass to (e) leave a remnant neutron star smaller than 1.5Mo.

(1648) Stars more massive than 20-30Mo may form (eb) a neutron star core, only to have it collapse due to (e) fallback of material from (e) the stellar mantle. Still more massive stars (above ~40Mo) may form (eb) a black hole directly in collapse, with or without (e&eb) a supernova explosion.

(1649) These masses are (=) rather uncertain at present, as our understanding of supernova explosion mechanisms is (=) still evolving [Woosley and Janka (2005), Muno (2006)].

(1650) Collapsing gas clouds larger than ~100Mo are expected to dissipate from (e) radiation pressure, which would prevent (e) more massive stars from forming, and consequently set (eb) an approximate upper limit for stellar mass black holes.

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| (1651) Observational evidence for stellar-mass black holes comes primarily from (e) X-ray astronomy. |
| (1652) A black hole in a close binary orbit can form (eb) an accretion disk (see Figure 1) of matter pulled off (er) a normal stellar companion. In its inspiral, disk material is heated by shearing, and becomes (=) a strong X-ray emitter. |
| (1653) If observations reveal a point-like X-ray emitter with (e&eb) a mass inferred from (e) orbital dynamics to be (=) above that possible for a neutron star, it becomes (=) a black hole candidate. |
| (1654) t present there are over a dozen such black hole candidates, including the first object to be identified as (=) a black hole, the X-ray source Cygnus X-1 (estimated mass 10±3Mo) [Casares (2006)]. |
| (1655) The other class of observationally-supported black holes is (=) that of supermassive black holes, ranging from hundreds to (e&eb) billions of solar masses. |
| (1656) Evidence for the existence of black holes at (eb) the upper end of this range is overwhelming. |
| (1657) By contrast, the existence of black holes with (e&*eb) mass roughly of order 102 to 104Mo, referred to as "intermediate-mass black holes," is speculative. |
| (1658) The supermassive holes may have begun as (=) primordial or stellar-mass black holes, but have grown through (e&eb) absorption of stars or gas or through mergers with (e&eb) other holes [Ferrarese and Ford (2005)]. |
| Figure 4: Observed orbits of stars near (eb) the dark compact object at the center of (e) our Milky Way galaxy: still frame from an animation. Animation created by Prof. Andrea Ghez and her research team at UCLA, from data sets obtained with the W. M. Keck Telescopes: http://www.astro.ucla.edu/~ghezgroup/gc/ |
| Evidence for supermassive black holes originally came from (e) quasars, small but intense radio sources seen at (eb) cosmological distances. |
| Their huge luminosities implied (eb) high mass, while (e) their rapid variability implied (eb) extremely small size. |
| More recently, measurements of Doppler shifts of stars and gas in the centers of galaxies have shown (eb) that compact objects of mass greater than 106Mo reside in (eb) the cores of most galaxies; the very small size of these central objects rules out any plausible alternative to the black hole explanation. |
| The black hole at the center of our own Milky Way galaxy has (e) a measured mass of ∼3×106Mo , based on (e) the orbits of stars near the radio source Sagittarius A* associated with (e&eb) the Galactic core [Ghez et al. (2005); see Fig. 4]. |
| (1664) Black holes are of particular interest for (e) gravitational waves, and vice versa. |
| (1665) For gravitationally-bound systems, the typical maximum gravitational wave amplitude (dimensionless strain) is (=) h∼(2GM/c2)2/(RD)∼R2H/(RD) , where D is the distance to the system, M and R are the mass and size of (e) the system, and RH is the horizon radius of (e) a black hole of the same mass. |
| (1666) From this it is clear that strongest gravitational waves will involve (e&eb) systems at or near black hole compactness: in particular, systems containing (er) black holes in close proximity to one another [Flanagan and Hughes (2005)]. |
| (1667) Conversely, a system containing only black holes would not be expected to radiate (eb) any sort of radiation other than gravitational waves, and in any system containing (e) black holes, gravitational waves provide (eb) the most accurate probes of conditions deep in (eb) their gravitational wells. |
| (1668) In particular, gravitational wave observations could definitively test (e&eb) whether a compact mass is truly a black hole [Cutler and Thorne (2002)]. |

(1669) The strongest intrinsic source of gravitational waves in the present Universe is (=) the collision of two comparable-massed black holes.

(1670) A sufficiently tight black hole binary system will lose (e) energy through emission of (e&eb) gravitational waves, causing (eb) the orbit to circularize and then slowly shrink over time.

(1671) Figure 5: Gravitational waves from (e) a black hole binary inspiral: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/

(1672) When the orbital radius approaches (eb) the Schwarzschild radius of the system, complex nonlinear dynamics come into play (e&eb) and the final stages of (e) the inspiral and merger can only be modeled through (e&eb) numerical simulations.

(1673) These simulations are quite challenging, but recent breakthroughs [Pretorius (2005)(See Fig. 5), Campanelli et al.(2006), Baker et al.(2006)] have led to (eb) the first complete inspiral-merger gravitational waveforms.

(1674) After the black holes have coalesced, they form (eb) a single highly-distorted black hole which quickly settles down to (e) a quiescent Kerr state through (e&eb) mission of quasinormal "ringdown" gravitational waves [Pretorius (2005), Campanelli et al. (2006), Baker et al. (2006)].

(1675) Another gravitational emission mechanism for supermassive black holes would come from (e) their capture of compact stellar objects, particularly stellar-mass black holes.

(1676) Unlike comparable-mass mergers, these captures would be on (eb) highly eccentric orbits that would not circularize before merger, resulting in (eb) intricate rapidly precessing orbits in (eb) the final months or years of inspiral.

**(1677)** The consequent gravitational-wave signal would allow (eb) for exquisitely precise measurements of the parameters of (e) the supermassive black hole, as well as tests of any deviation from (e) the predictions of general relativity [Cutler and Thorne (2002)].

**(1678)** For stellar-mass black holes (tens of Mo), the final stages of inspiral and merger occur at (eb) frequencies of hundreds to thousands of Hz, a regime targeted by (e) ground-based gravitational-wave observatories.

**(1679)** Systems of supermassive black holes that are (=) millions of Mo emit at frequencies of tens of (e) mHz, which will be detectable by (e) proposed space-based (e) gravitational-wave detectors.

**(1680)** Black holes in the billions of Mo, or a stochastic background of signals from (e) more moderate supermassive black holes at (eb) earlier stages of their inspiral, might emit (eb) waves at nHz frequencies that would be measurable through (e&eb) years-long correlated timing measurements of radio pulsars [Jenet et al. (2005)].

**An equation to quantify the origins of life on other planets July 5, 2016 by Bob Yirka report**

(1681) (Phys.org)—A pair of researchers, one with the Columbia Astrobiology Center in New York, the other with the University of Glasgow in the U.K. has come up with a mathematical equation that when (e) solved is meant to offer (eb) a means for estimating how often life begins on other planets. In their paper published in Proceedings of the National Academy of Sciences, Caleb Scharf and Leroy Cronin describe their equation, how they came up with it and why they believe it might become more useful as scientists learn more about the true nature of other planets and solar systems.

(1682) Back in the 1960's, an astronomer named Frank Drake came up with a formula for (e) estimating the likely number of (e) alien civilizations that might be capable of transmitting radio signals in such a way as to be (=) recognizable by receivers here on earth.

(1683) That equation, named quite naturally the **Drake Equation** has been the unofficial standard-bearer for half a century, despite it having no solution because (e) some of the parameters are still unknown.

(1684) The same holds true for the new equation developed by (e) Scharf and Cronin, but the results given should the unknowns ever be discovered would offer a different sort of answer than Drake was looking for.

(1685) To build their formula, the duo came up with several parameters to solve for (e) **abiogenesis**, which is the likelihood of an event occurring that leads to (eb) life beginning, they included (e) the number of possible building blocks, the mean number of such blocks per possible organism, the availability of building blocks that might exist during a given time period, expressed as a fraction and the probability that the existence of the building blocks would actually lead to (eb) life starting in a given unit of time. It looks like this: **Nabiogenesis (t)= Nb · 1/no · fc · Pa · t**

(1686) In essence the formula is meant to suggest (eb) that the probability of life beginning on a given planet is very likely connected to (e&eb) whether there are building blocks available on a given planet, and how much of them there might be.

(1687) As more is learned by space researchers it is hoped that the formula could help narrow down (e) search targets by offering (eb) a statistical probability of success for (e) a given planet.

(1688) The researchers note that the building blocks available do not necessarily have to be (=) of the type that led to (eb) life beginning on our planet.

**(1689)** They also suggest that implied in the formula is (=) the likelihood that life beginning events are (=) more likely to occur in solar systems where (e) there are multiple planets, which allow (eb) for an opportunity to share materials that could lead to (eb)building blocks **More information: Quantifying the origins of life on a planetary scale, Proceedings of the National Academy of Sciences, www.pnas.org/cgi/doi/10.1073/pnas.1523233113**

**Abstract**

(1690) A simple, heuristic formula with (e&eb) parallels to the Drake Equation is introduced to help focus discussion on open questions for (e) the origins of life in a planetary context.

(1691) This approach indicates (eb) a number of areas where (e) quantitative progress can be made on parameter estimation for (e) determining origins of life probabilities, based on (e) constraints from Bayesian approaches.

(1692) Authors discuss a variety of "microscale" factors and their role in determining (eb) "macroscale" abiogenesis probabilities on (e&eb) suitable planets.

(1693) Authors also propose that impact ejecta exchange between (e&eb) planets with parallel chemistries and (e&eb) chemical evolution could in principle amplify (eb) the development of (e) molecular complexity and (e&eb) abiogenesis probabilities.

(1694) This amplification could be very significant, and both bias conclusions about abiogenesis probabilities based on (e) the Earth and provide (eb) a major source of variance in the probability of life arising in (eb) planetary systems.

**(1695)** Authors use heuristic formula to suggest (eb) a number of observational routes for improving constraints on (e) origins of life probabilities. **Read more at: http://phys.org/news/2016-07-equation-quantify-life-planets.html#jCp**

**Brian Koberlein Shared publicly - Yesterday 8:00 PM the Illusion of Truth**

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| (1696) | Physical theories are often presented as (=) a description of what's really going on. |
| (1697) | Forces act on (e&eb) a baseball, causing (eb) baseball to fall. |
| (1698) | Atoms collide and fuse in (eb) the heart of a star, releasing heat and energy. Science is true, as is often said, and our scientific theories encapsulate this truth. But this isn't entirely true. |
| (1699) | For example, quantum mechanics is a strange, sometimes confusing theory. Objects can be (=) particles and waves. |
| (1700) | They do strange things like tunnel through (e&eb) barriers, and can even appear and disappear (e&eb). |
| (1701) | Given all this strangeness, what's really going on? What's the true nature of quantum reality? It depends on (=) which approach you want to take. |
| (1702) | In one view quantum objects are described by (e) a wave-like probability function. |
| (1703) | When you interact with (e&eb) or "observe" (e) the object this "wavefunction" collapses into (eb) a definite state. |
| (1704) | That's view commonly presented, but a quantum system can also be described by (e) its transitions between energy states. |
| (1705) | Since this uses a mathematical method involving (e&eb) matrices, it's known as (=) matrix mechanics. |
| (1706) | Both wavefunctions and matrix mechanics give (eb) the same results, but their view of what really goes on is very different. |
| (1707) | Then there's the path integral method |
| (1708) | . Rather than a wavefunction or matrix transitions, path integrals imagine quantum objects can take almost any path between (e&eb) two states. |
| (1709) | By summing all the possible paths you can derive (eb) the odds that it will occur. |
| (1710) | So which one is true? Are quantum objects distributed waves of (e) probability? |
| (1711) | Are they simply transitions between (e&eb) energy states? |
| (1712) | Do they take an infinite number of (e) paths between states? |
| (1713) | Each of these models makes (eb) the same predictions, so one version is no more "true" than the others. |
| (1714) | What happens in practice is that we'll use whatever method is useful at the time. They are equivalent models, so the best model for the job is the one we'll use. The only reason the wavefunction view is so common is that it's (=) the version usually taught to introductory students. You might think this uncertainty of truth is due to (e) the behavior of quantum physics itself. It's so strange and counterintuitive that we can't (e) wrap our puny brains around what's really going on. But the same thing occurs in lots of other fields. Even something as straight forward as basic Newtonian physics |
| (1715) | Toss a baseball in the air and (eb) Earth's gravitational force pulls it down. |
| (1716) | The force of gravity is a simple truth, right? While we often describe classical motion in terms of (e&eb) forces and acceleration, we can also describe it in terms of (e&eb) energy and momentum. |
| (1717) | In the Lagrangian and Hamiltonian approach, the path of a baseball is (=) the optimized path among possibilities. |
| (1718) | In this view a baseball's path is (=) the extrema of an energy equation, and force can be derived as a necessary consequence of (e) this. |
| (1719) | In the relativistic view the baseball follows (e&eb) a geodesic, which is (=) the minimal path through (e&eb) space and time. |

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| (1720) So is a baseball's motion due to (e) a gravitational force, an energy extrema, or a spacetime geodesic? |
| (1721) As with quantum theory, different approaches yield (eb) the same result. They are mathematically equivalent, so we can use whatever method is most useful at the time. |
| (1722) At its core, science is less about truth and more about models. The metaphysics underlying a model is useful only as far as it allows us (eb) to make better predictions, generate (eb) new ideas, or bring models together as (=) a cohesive whole |
| (1723) This is why we have no problem using (e) classical gravity to calculate the path of (e) a spacecraft through (e&eb) the solar system, while (e) using special relativity to account for (e) the Doppler shift of the spacecraft's radio signals. |
| (1724) It's why we can use (e) quantum physics to study atoms in the morning, and general relativity to (e) study black holes in the afternoon. |
| (1725) In regimes where models conflict with each other it isn't a failure of truth, but instead shows an opportunity to develop a better model. |
| (1726) It could be that with each better model we move closer to (e&eb) the truth about reality. |
| (1727) The truth is out there, and science strives to move towards (e&eb) that truth. |
| (1728) It's a common view, and certainly the search for truth has driven (eb) many scientists to develop better and better models. |
| (1729) But the real power of science is the recognition that what we have are models. Our models can be powerful, but they are always a bit tentative. There's always a chance that they might just be an illusion of truth. |

**Synopsis: A Single-Photon Cheshire Cat July 1, 2016 Researchers detected the polarization of a photon separate from the photon itself, just as the grin of Lewis Carroll's Cheshire cat can appear apart from the cat's body Synopsis figure APS/Carin Cain**

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| (1730) In a quantum Cheshire cat, one property of a particle is present detached from (e) the particle itself, just as (=) the grin of the feline in Alice in Wonderland can exist in the absence of (e) the whole animal. |
| (1731) Now a research team has successfully separated (e&eb) a single photon from (e) its polarization. |
| (1732) The effect was previously demonstrated (eb) with neutrons and also with light in a classical scenario, but the latest experiment uses single photons. |
| (1733) To produce the Cheshire cat, **Maximilian Schlosshauer and his colleagues** at the University of Portland in Oregon used an interferometer where (e) light hits a beam splitter, goes in two, opposite directions around a circuit (but along slightly offset routes), and then recombines at (eb) the entrance. |
| (1734) The incoming polarization of light was horizontal, but this was switched to (e&eb) vertical for the counterclockwise loop. |
| (1735) Each path included a device to adjust (e&eb) the polarization further and another to (e&eb) scatter away a small number of photons of (e) a designated polarization, the equivalent of detecting their presence. |
| (1736) The team fired single photons into (e&eb) the interferometer so that only one photon was present at (eb) a time, traversing both paths simultaneously. |
| (1737) They placed a horizontal polarizer in front of the output photon detector and gradually varied (e&eb) the difference between the two path lengths. |

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| (1738) This variation generated (eb) a sine-wave interference pattern from (e) the exiting photons only when the polarization was tweaked in (eb) the counterclockwise path. |
| (1739) At the same time, the average number of photons counted decreased only when (e) light was scattered from the clockwise path. |
| (1740) Simultaneously demonstrating these different effects for two routes of a single photon constitutes (eb) evidence of the Cheshire cat. |
| **(1741)** The effect might someday improve precision experiments if (e) an unwanted property could be moved away from the property being measured. **This research is published in Physical Review A. – David Ehrenstein David Ehrenstein is the Focus Editor for Physics.** |

**Observation of the quantum paradox of separation of a single photon from one of its properties James M. Ashby, Peter D. Schwarz, and Maximilian Schlosshauer Phys. Rev. A 94, 012102 – Published 1 July 2016**

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| (1742) Authors report an experimental realization of the quantum paradox of (e) the separation of a single photon from (e) one of its properties (the so-called "quantum Cheshire cat"). |
| (1743) Authors use a modified Sagnac interferometer with (e&eb) displaced paths to produce (eb) appropriately pre- and postselected states of (e) heralded single photons. |
| (1744) Weak measurements of photon presence and (e&eb) circular polarization are performed in each arm of (e) the interferometer by introducing (eb) weak absorbers and small polarization rotations and analyzing changes in (eb) the postselected signal. |
| (1745) The absorber is found to have (e) an appreciable effect only in one arm of the interferometer, while (e) the polarization rotation significantly affects (e&eb) the signal only when (e) performed in the other arm. |
| (1746) Authors carry out both sequential and simultaneous weak measurements and find (eb) good agreement between measured and predicted weak values. |
| (1747) In the language of Aharonov et al. and in the sense of the ensemble averages described by (e) weak values, the experiment establishes (eb) the separation of a particle from one its properties during the passage through (e&eb) the interferometer. |

**Viewpoint: Relativity Gets Thorough Vetting from LIGO Frans Pretorius, Department of Physics, Princeton University, Princeton, NJ 08544, USA May 31, 2016• Physics 9, 52 The gravitational-wave signal observed by the LIGO detectors shows no deviation from what general relativity predicts. Figure caption expand figure Frans Pretorius, APS/Carin Cain**

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| (1748) Figure 1: The signal from one of the LIGO detectors in Hanford, Washington, is shown with two representations of (e) the best-fit numerical relativity (NR) waveform. |
| (1749) The filtered NR waveform illustrates how the raw waveform is perceived by (e) the detector, |
| (1750) On September 14, 2015, the LIGO interferometers made the first direct measurement of (e) gravitational waves [1]. |
| (1751) The historic discovery came nearly 100 years after Einstein showed that his general relativity theory predicts (eb) the existence of gravitational waves and over 50 years since (e) the first quest to observe the Universe in gravitational waves [2]. |
| (1752) The observed event, labeled GW150914, was (=) the collision of two black holes, one weighing 29 solar masses, the other 36 solar masses. |

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| (1753) This binary coalesced into (e&eb) a remnant black hole weighing 62 solar masses and rotating (e&eb) at 67% the maximum rate allowed (eb) for a black hole. |
| (1754) The most statistically significant portion of (e) the signal came from (e) the last several gravitational-wave cycles of the collision, where (e) the most nonlinear, strong-field dynamics of the theory manifest. |
| (1755) One couldn't have asked for a better first event to learn about (e&eb) this regime of general relativity. |
| (1756) The first thing that GW150914 can teach us is how closely nature follows (eb) Einstein's prescriptions for (e&eb) colliding black holes. |
| (1757) Researchers from the LIGO team and the Virgo team have performed a series of tests to see if (e) any violations of (e) general relativity might be lurking in the data of GW150914 [3]. |
| (1758) They first determined (eb) the theoretical waveform, or "template," from (e) general relativity simulations that best fit the measured signal (see Fig. 1). |
| (1759) Their main finding is that the residual signal—obtained by (e) subtracting this template from (e) the data—is consistent with (e&eb) noise. |
| (1760) This is, in a sense, the "everything" test, as (e) the template folds in (e&eb) all that general relativity predicts about the merger: the decay of the last few orbits of (e) the two black holes, their violent collision (e&eb) and astonishingly rapid relaxation to (e) a quiescent rotating black hole, the propagation of (e) gravitational waves through (e&eb) the cosmos, and finally the transformation of (e) these waves into(e&eb) a chirp in the LIGO detectors. |
| (1761) Given the large signal-to-noise ratio of 24, the close agreement between the template and (=) the data puts limits on (e) alternative models. |
| (1762) Specifically, the predictions of these models can only disagree with (e&eb) the general relativity prediction for GW150914 by, at most, 4%. |
| (1763) In addition to this comprehensive test, the LIGO/Virgo teams also performed other checks that pick apart the signal, in order to address particular aspects of general relativity and black hole physics. In one such test, they compared the early and (e&eb) late parts of the waveform. |
| (1764) From the waves generated during (e&eb) the initial inspiral, one can infer (eb) the masses and spins of the binary. |
| (1765) Given these numbers, general relativity allows (eb) the mass and spin of (e) the remnant black hole to be calculated. |
| (1766) The researchers compared this result to (e&eb) an independent measurement of (e) the mass and spin obtained from (e) the late-time ringdown, and to within (eb) uncertainties, the two estimates were (=) consistent. |
| (1767) Another test isolated coefficients in the so-called Post Newtonian (PN) series describing (eb) the inspiral. |
| (1768) The PN series is (=) an expansion of the gravitational-wave signal in powers of (e) the velocity vv of the sources relative to (e&eb) the speed of light cc. |
| (1769) Observations of binary pulsars, which show (eb) a decrease in their orbital period due to (e) gravitational-wave emission, have constrained (e) PN terms, especially at low order. |
| (1770) For higher-order terms associated with (e&eb) strong-field emission, the LIGO/Virgo teams were able to improve (eb) the limits by many orders of magnitude compared to (e&eb) the binary pulsar constraints. |

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| (1771) At first glance it might seem these new results simply add to (e) the litany of previous tests that general relativity passed with flying colors. |
| (1772) However, GW150914 is vastly more significant than that. Other tests have shown that general relativity is valid in (eb)the weak-field regime, but GW150914 gives (eb) the first direct observational evidence that general relativity is the correct theory of nature governing (e&eb) the extreme spacetime dynamics around the event horizon of (e) a black hole. |
| (1773) One could argue that binary pulsar observations already told us that gravitational waves exist [4], but GW150914 is more luminous in gravitational waves than pulsar binaries by (e) at least 24 orders of magnitude. |
| (1774) It's true that the evolution of the Universe as a whole is driven by (e) strong gravity, but it would be difficult to argue that precision cosmological measurements can constitute (eb) a test of general relativity when (e) one of the dominant inputs, dark energy, is still a mystery. |
| (1775) As far as black holes are concerned, astronomers have (e) observed ultra compact dark objects that are consistent with (e&eb) black holes. |
| (1776) The most striking data comes from (e) stellar orbits about the 4 million solar mass objects in (eb) the center of our Galaxy [5], but present observations cannot constrain it to be (=) a black hole as described by general relativity. |
| (1777) GW150914 has given us the first morsels of quantifiable evidence that the rich physics of black holes is as predicted by general relativity. For example, the LIGO observations are consistent with (e&eb) so-called cosmic censorship [6], which says that gravitational singularities must remain hidden behind an (eb) event horizon. |
| (1778) If a "naked" singularity were produced in (eb) the energetic collision of GW150914, the data would likely have shown (eb) some sort of anomaly, but in exactly what form, we don't know. |
| (1779) Another example involves the "no-hair" theorem [7], which says that the only properties of a black hole are (=) its mass, spin, and charge (though astrophysical black holes are expected to have negligible charge). |
| (1780) By consequence, these properties should determine (eb) the frequency and decay constants of (e) all the infinitely many ringdown modes of (e) a black hole. |
| (1781) With GW150914, only the least-damped ringdown mode is (=) visible above the noise, but when combined with (e&eb) the remnant mass and spin inferred from (e) the inspiral, this is enough to count as (=) the first piece of evidence for a black hole's baldness. |
| (1782) GW150914 should also be a test for (e) alternative theories of gravity or exotic alternatives to (e) black holes within general relativity, such as boson stars [8]. |
| (1783) Some of these alternative models are concerned with (e&eb) explaining dark energy or dark matter and do not (e) therefore, predict any deviation from relativity at (eb) the scale of GW150914. |
| (1784) However, other models do in fact deviate from (e) conventional relativity in their predictions of (e) the gravitational-wave signal from (e) a binary merger. |
| (1785) The largest deviations should presumably occur during (e&eb) the last several gravitational-wave cycles, where (e) the signal is strong, and deviations should be (=) easy to detect or constrain. |
| (1786) Embarrassingly though, the majority of alternatives that claim deviations can't (e) really say much about this event, as the two-body problem hasn't been solved in these theories. |
| (1787) These models can only reproduce (eb) a fraction of the full gravitational-wave signal, (eb) so the strong constraints one might have hoped for (e) cannot be leveled. |

(1788) Exotic alternatives within (eb) general relativity are in a similar boat, or worse for (e) objects like gravastars, which have (e) no well-defined theoretical framework to say what happens when (e) two such objects collide.

(1789) Regardless, what is clear from the data is that any non-black-hole solution must imagine a form of (e) exotic matter with a viscosity larger than (e) any known material to explain (eb) the rapid damping of oscillations observed during the ringdown.

(1790) The inferred event rate for heavy black hole mergers like GW150194 is high, and LIGO has another factor of 3 improvement in sensitivity slated before it reaches (eb) design sensitivity.

(1791) Within 5 years then, we can expect dozens of detected events, with the loudest having a signal-to-noise ratio upwards of (e) 100.

(1792) This will allow (eb) for exquisite insight into the nature of extreme gravity beyond (e) what GW150194 has already gifted us with. **This research is published in Physical Review Letters**.

**Tests of General Relativity with GW150914 B. P. Abbott et al. (LIGO Scientific and Virgo Collaborations) Phys. Rev. Lett 116, 221101 – Published 31 May 2016**

(1793) The LIGO detection of GW150914 provides (eb) an unprecedented opportunity to study the two-body motion of (e) a compact-object binary in (eb) the large-velocity, highly nonlinear regime, and to witness the final merger of (e) the binary and the excitation of uniquely relativistic modes of (e) the gravitational field.

(1794) Authors carry out several investigations to determine whether GW150914 is consistent with (e&eb) a binary black-hole merger in general relativity.

(1795) They find that the final remnant's mass and spin, as determined from (e) the low-frequency (inspiral) and (e&eb) high-frequency (postinspiral) phases of the signal, are mutually consistent with (e&eb) the binary black-hole solution in general relativity.

(1796) Furthermore, the data following the peak of GW150914 are consistent with (e&eb) the least-damped quasinormal mode inferred from (e) the mass and spin of the remnant black hole.

(1797) By using waveform models that allow (eb) for parametrized general-relativity violations during (e&eb) the inspiral and merger phases, authors perform quantitative tests on (e&eb) the gravitational-wave phase in the dynamical regime and they determine the first empirical bounds on (e&eb) several high-order post-Newtonian coefficients.

(1798) Authors constrain (e) the graviton Compton wavelength, assuming (eb) that gravitons are dispersed in (eb) vacuum in the same way as particles with mass, obtaining a 90%-confidence lower bound of 1013 km.

(1799) In conclusion, within our statistical uncertainties, we find no evidence for violations of (e) general relativity in the genuinely strong-field regime of gravity. **Received 26 March 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.221101 © 2016 American Physical Society**

**Focus: Background Noise of Gravitational Waves March 31, 2016• Physics 9, 33**

(1800) Assuming LIGO's recently detected black hole merger is not unusual; researchers revised upward their estimate of (e) the strength of the background noise coming from (e) distant mergers across the Universe. Figure caption expand figure Caltech/MIT/LIGO Lab (deleted due to spatial constraints)

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| (1801) Catching some waves The LIGO-Livingston interferometer consists of (e) two 4-km-long arms where (e) laser light bounces back and forth. Interference between the lasers from the two arms can signal tiny changes in (e&eb) the lengths of the arms resulting from gravi... |
| (1802) The recent detection of gravitational waves emitted from (e) a pair of merging black holes implies (eb) the existence of a large population of similar pairs, which should be generating (eb) a background of waves too faint to be detected with (e&eb) present technology. |
| (1803) Using this information, researchers now estimate that this background could be (=) roughly 10 times stronger than previously expected and might be found by (e) sensitive instruments in only a few years. |
| (1804) The background signal could help astrophysicists better understand how black hole binaries form and also pave the way to detecting the waves emitted by (e) the expanding Universe shortly after the big bang. |
| (1805) Gravitational waves were detected for the first time on 14 September 2015 by the Laser Interferometer Gravitational-Wave Observatory (LIGO). |
| (1806) By reflecting laser light between mirrors set 4 km apart, the team was able to detect (e&eb) tiny fluctuations in the distance between the mirrors as (=) the waves passed. |
| (1807) The waves were triggered by (e) an event, called GW150914; involving (e&eb) the merger of a pair of black holes having (e) masses of about 29 and 36 times that of the Sun. |
| (1808) Researchers expect that similar pairs of black holes, as well as neutron stars, should be merging constantly, creating (eb) gravitational waves that should bombard (e&eb) the Earth from all directions and make the equivalent of (=) a faint popping noise in a detector. |
| (1809) Previous estimates of the likely strength of this background had to make assumptions about (e&eb) how many pairs there might be and about the distribution of their masses. |
| (1810) These estimates remained speculative, as no actual binaries were known. |
| (1811) “What the recent observation has changed,” says Nelson Christensen of Carleton College in Minnesota and the LIGO Scientific Collaboration (LSC), “is that we now know there really is (=) a population of black hole binaries. We have a small bit of information about it.” |
| (1812) The new analysis—carried out by LSC researchers working with (e&eb) the European Virgo collaboration—estimates (eb) the energy density in the gravitational wave background in (eb) the frequency range below 100 Hz, where (e) the LIGO detector is most sensitive. |
| (1813) To make this estimate, they assumed (eb) a population of black hole binaries of various sizes, having (e) a range of distances from Earth, and created through (e&eb) a variety of scenarios. |
| (1814) They then calculated the amount of gravitational radiation that would result from (e) the merging of these objects. |
| (1815) The team then adjusted this estimate by using (e) details of the GW150914 event—especially the masses of the black holes involved—to make (eb) a most likely guess at the actual distribution of black hole binaries. |
| (1816) The idea is that the observed merger is (=) most likely representative of (e) the general population, rather than being extremely unusual. |
| (1817) If the black hole binary population is represented as (=) a bell curve, or Gaussian, the event should lie within (eb) the main bulge, not in the tails. |
| (1818) The observed mass of GW150914 suggests (eb) that the main population of black hole binaries could be (=) roughly 20 times more massive than expected, which would generate (eb) a background roughly 10 times “louder” than (e) previous estimates. |

(1819) However, Christensen emphasizes that the result has (e) lots of uncertainty, as it's based on (e) only one event.

(1820) "This estimate seems very plausible to me," says astrophysicist Maura McLaughlin of the University of West Virginia in Morgantown. She is a member of the NANOGRAV collaboration, which tries to detect gravitational waves at much lower frequencies than LIGO can observe. "I can't see any holes in their arguments. Unless LIGO got really lucky, then there must be many more of these systems than previously thought." This level of gravitational background implies (eb) that the LIGO and Virgo detectors might be sensitive enough to see it in 2–5 years, sooner than physicists previously expected, says Christensen, which has led to much anticipation. "Everyone is super excited," he says.

(1821) Detecting the background would provide (eb) a confirmation of the presence of gravitational waves carrying (e&eb) energy throughout the Universe. It would also be useful, the researchers suggest, in gaining (e) further understanding of precisely how black holes form.

(1822) And subtracting this background might someday allow (eb) detection of the fainter waves coming from (e) the expansion of the Universe moments after the big bang. This research is published in Physical Review Letters. –Mark Buchanan Mark Buchanan is a freelance science writer in Normandy, France.

**GW150914: Implications for the Stochastic Gravitational-Wave Background from Binary Black Holes B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 131102 – Published 31 March 2016**

(1823) The LIGO detection of the gravitational wave transient GW150914, from (e) the inspiral and merger of (e) two black holes with masses $\gtrsim 30 M_{\odot}$, suggests (eb) a population of binary black holes with (e&eb) relatively high mass.

(1824) This observation implies (eb) that the stochastic gravitational-wave background from binary black holes, created from (e) the incoherent superposition of (e) all the merging binaries in the Universe, could be (=) higher than previously expected.

(1825) Using the properties of GW150914, authors estimate the energy density of (e) such a background from (e) binary black holes.

(1826) In the most sensitive part of the Advanced LIGO and Advanced Virgo band for (e) stochastic backgrounds (near 25 Hz), authors predict $\Omega_{GW}(f=25\ \mathrm{Hz}) = 1.1^{+2.7}_{-0.9} \times 10^{-9}$ with 90% confidence.

(1827) This prediction is robustly demonstrated for (e) a variety of formation scenarios with (e&eb) different parameters.

(1828) The differences between models are (=) small compared to the statistical uncertainty arising from (e) the currently poorly constrained (e) local coalescence rate.

(1829) Authors conclude that this background is potentially measurable by (e) the Advanced LIGO and Advanced Virgo detectors operating at their projected final sensitivity. DOI:http://dx.doi.org/10.1103/PhysRevLett.116.131102 © 2016 American Physical Society

**Study opens new realms of light-matter interaction July 15, 2016 by David L. Chandler. Some sentences might have been omitted due to spatial constraints. Please bear with me.**

(1830) Emission spectra are a widely used method for (e) identifying chemical compounds; the bright lines reveal (eb) the different frequencies of light that can be emitted by (e) an atom.

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| (1831) Here, a normal emission spectrum for (e) an atom in a high-energy state (top) is compared to (e&eb) the emission from the same atom placed just a few nanometers (billionths of a meter) away from (e) graphene that has been doped with (e&eb) charge carriers (bottom). |
| (1832) For each energy-level transition, an orange line (or purple cloud) appears if (e) that transition is estimated to be faster than (e) one per microsecond — making it frequent enough to be (=) observed. **Credit: Massachusetts Institute of Technology** |
| (1833) A new MIT study could open up new areas of technology based on (e) types of light emission that had been thought to be (=) "forbidden," or at least so unlikely as to be (=) practically unattainable. |
| (1834) The new approach, the researchers say, could cause (eb) certain kinds of interactions between light and (e&eb) matter, which would normally take billions of years to happen, to take place instead (e&eb) within billionths of a second, under (e&eb) certain special conditions. The findings, based on a theoretical analysis, are reported today in the journal Science in a paper by MIT doctoral student Nicholas Rivera, Department of Physics Professor Marin Soljačić, Francis Wright Davis Professor of Physics John Joannopoulos, and postdocs Ido Kaminer and Bo Zhen. |
| (1835) Interactions between light and matter, described by (e) the laws of quantum electrodynamics, are the basis of a (e) wide range of technologies, including (e) lasers, LEDs, and atomic clocks. |
| (1836) But from a theoretical standpoint, "Most light-matter interaction processes are 'forbidden' by (e) electronic selection rules, which limits (e) the number of transitions between energy levels we have access to," Soljačić explains. |
| (1837) For example, spectrograms, which are used to (e) analyze the elemental composition of materials, show (eb) a few bright lines against (e&eb) a mostly dark background. |
| (1838) The bright lines represent (eb) the specific "allowed" energy level transitions in (eb) the atoms of that element that can be accompanied by (e) the release of a photon (a particle of light). |
| (1839) In the dark regions, which make up most of the spectrum, emission at those energy levels is (=) "forbidden." |
| (1840) With this new study, Kaminer says, "we demonstrate theoretically that these constraints can be lifted" using (e) confined waves within (e&eb) atomically thin, 2-D materials. |
| (1841) "We show that some of the transitions which normally take the age of the universe to happen could be made to happen within (eb) nanoseconds. |
| (1842) Because of this, many of the dark regions of a spectrogram become (=) bright once an atom is (=) placed near a 2-D material." |
| (1843) Electrons in an atom have (e) discrete energy levels, and when they hop from one level to another they give off (eb) a photon of light, a process called spontaneous emission. |
| (1844) But the atom itself is much smaller than the wavelength of the light that gets emitted—about 1/1,000 to 1/10,000 as big—substantially impairing (e&eb) the interactions between the two. |
| (1845) The trick is, in effect, to "shrink" (e) the light so it better matches the scale of the atom, as (=) the researchers show in their study. |
| (1846) The key to enabling a whole range of interactions, specifically transitions in (eb) atomic states that relate to absorbing or emitting light, is (=) the use of a two-dimensional material called graphene, in which light can interact with (e&eb) matter in the form of plasmons, a type of electromagnetic oscillation in the material. |
| (1847) These plasmons, which resemble photons but have wavelengths hundreds of times shorter, are (=) very narrowly confined in the graphene, in a way that makes some kinds of interactions with (e&eb) that matter many orders of magnitude more likely than (e) they would be in ordinary materials. |

(1848) This enables a variety of phenomena normally considered unattainable, such as (e) the simultaneous emission of multiple plasmons, or two-step light-emitting transitions between (e&eb) energy levels, the team says.

(1849) This method can enable the simultaneous emission of (e) two photons that are "entangled," meaning (eb) they share the same quantum state even when (e) separated.

(1850) Such generation of entangled photons is (=) an important element in quantum devices, such as those that might be used for cryptography.

(1851) Making use of these forbidden transitions could open up (eb) the ability to tailor the optical properties of materials in ways that had not been thought possible, Rivera says.

(1852) "By altering these rules" about the relationship between light and (e&eb) matter, "it can open (eb) new doors to reshaping (e&eb) the optical properties of materials."

(1853) Kaminer predicts that this work "will serve as (=) a founding piece for (e) the next generation of studies on light-matter interactions" and could lead to (eb) "further theoretical and experimental advances in many fields which rely on (e) light-matter interactions, including (e) atomic, molecular and optical physics, photonics, chemistry, optoelectronics, and many others."

(1854) Beyond its scientific implications, he says, "this study has possible applications across multiple disciplines, since (e) in principle it has potential to enable the full use of (e) the periodic table for optical applications."

(1855) This could potentially lead to (eb) applications in spectroscopy and sensing devices, ultrathin solar cells, and new kinds of materials to absorb solar energy, organic LEDs with higher efficiencies, and photon sources for (e) possible quantum computing devices."From the standpoint of fundamental science, this work lays the groundwork for a subfield that just a few years ago was difficult to imagine and until now was largely unexplored," Soljačić says.

(1856) "Two-dimensional materials confine (e) fields to a surface and motion to a plane, making possible (eb) many effects that are orders-of-magnitude too weak to (e) appear in a bulk volume," says Jason Fleischer, an associate professor of electrical engineering at Princeton University, who was not involved in this research.

(1857) This work, he says, "systematically explores how 2-D materials improve light-matter interactions, laying a theoretical foundation for faster electronic transitions, enhanced sensing, and better emission, including the compact generation of broadband and quantum light." More information: N. Rivera et al: Shrinking light to allow (eb) forbidden transitions on the atomic scale, Science (2016). DOI: 10.1126/science.aaf6308 Journal reference: Science search and more info website: Provided by: Massachusetts Institute of Technology search and more info website. This story is republished courtesy of MIT News (web.mit.edu/news office/), a popular site that covers news about MIT research, innovation and teaching.

**Viewpoint: Peering Back to the Beginning of Time Lawrence M. Krauss, Department of Physics and School of Earth and Space Exploration, Arizona State University, Tempe, AZ 85287, USA and Mount Stromlo Observatory, Research School of Astronomy and Astrophysics, Australian National University, Weston, ACT 2611, Australia June 19, 2014• Physics 7, 64**

(1858) The BICEP2 collaboration reports the detection of BB-mode polarization of (e) the cosmic microwave background—a signal that might originate from (e) gravitational waves created by (e) inflation during the very earliest moments in the evolution of our Universe. Figure caption expand figure P. A. R. Ade et al. [1]; Image on homepage: BICEP2 Collaboration (deleted due to restrictions on space)

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| (1859) Figure 1: (Top) BB-mode polarization pattern observed by (e) the BICEP2 telescope [1]. |
| (1860) The twisting pattern in the polarization of (e) the cosmic microwave background (CMB) might originate from (e) gravitational waves created by inflation in the very first moments of the universe. |
| (1861) Inflation—the hypothesis that the Universe underwent (e&eb) a phase of superluminal expansion in a brief period following the big bang—has (e) the potential of explaining, from first principles, why the Universe has (e) the structure we see today. |
| (1862) It could also solve outstanding puzzles of standard big-bang cosmology, such as why the Universe is, to a very good approximation, (=) flat and isotropic (i.e., it looks the same in all directions). |
| (1863) Yet we do not yet have a compelling model, based on (e) fundamental particle physics principles, that explains inflation. |
| (1864) And despite its explanatory power and a great deal of suggestive evidence, we still lack an unambiguous and direct probe of (e) inflation. |
| (1865) Theorists have developed different models for inflation, which all share (e&eb) a common, robust prediction: Inflation would have created (eb) a background of gravitational waves that could have (eb) an observable effect. |
| (1866) These waves would cause (eb) subtle, characteristic distortions of (e) the cosmic microwave background (CMB)—the oldest light in the Universe, released when (e) photons decoupled from matter and the Universe became (=) transparent to radiation |
| (1867) Until now, such a background of gravitational waves has (e) escaped detection. It is therefore very exciting that the BICEP2 (Background Imaging of Cosmic Extragalactic Polarization) collaboration has now reported in Physical Review Letters the observation of polarization modes of CMB radiation [1] that are consistent with (e&eb) such a background. |
| (1868) If confirmed, BICEP2's finding would provide (eb) the first direct probe of physical processes that occurred a mere 10−3610-36 seconds after (e) the big bang. |
| (1869) The result also offers a rare opportunity to directly test theoretical ideas, including inflation and grand unification (the theory that merges the electromagnetic, weak, and strong forces into a single force), at energies 13 orders of magnitude greater than (e) we can currently access at the Large Hadron Collider. |
| (1870) Inflation is based on (e) the possibility that a scalar field in the early Universe got momentarily stuck in (eb) a "false vacuum state"—a metastable condition in which the Universe wasn't at (eb) its global potential minimum. |
| (1871) The energy stored in this vacuum state could drive (eb) an exponential expansion of spacetime (similar to, but much faster than, the expansion due to dark energy seen today) before (e) a phase transition to the true vacuum state occurs, releasing (eb) this energy into (e&eb) particles and radiation. |
| (1872) This creates (eb) an initially hot state for (e) a subsequent big-bang-like expansion from (e) which our Universe subsequently evolved [2]. |
| (1873) If the Universe grew by at least a factor of 10261026, (eb) inflation can explain (eb) both the apparent flatness and isotropy of the Universe [2]: |
| (1874) The expansion would stretch (eb+ eb-dot) any initial curvature of (e) the Universe to near flatness, and smooth out (e&eb) any initial inhomogeneities, just as stretching (eb+ ebdot) a piece of fabric would do. |
| (1875) Isotropy is justified by the fact that the Universe, before inflation, was once much smaller and could have been contained in (eb) a causally connected region that was in thermal equilibrium, resulting in (eb) the approximately uniform, isotropic temperature distribution observed on (eb) large scales today. |

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| (1876) Importantly, the theory can explain the (eb) origin of structure in the Universe: Microscopic quantum fluctuations associated with (e&eb) the scalar field driving inflation could have been "frozen in" during the exponential expansion, resulting, (eb) after inflation, in density fluctuations that might have seeded all observed large-scale structures in the Universe, including (e) all visible galaxies and clusters of galaxies. |
| (1877) The signs of such primordial fluctuations have indeed been detected through (e&eb) their effect on (e&eb) the CMB, first by the Cosmic Background Explorer (COBE) [3] and later by other missions. |
| (1878) These observational studies detected small temperature variations (anisotropies) in (eb) the CMB and found a spectrum compatible with (e&eb) that predicted by inflation. |
| (1879) The magnitude of these fluctuations suggests (eb) an energy scale of inflation (loosely speaking, the energy stored throughout space in (eb) the scalar field potential that drove inflation) close to (e) 10161016 giga-electron-volts (GeV). |
| (1880) But while consistent with inflation, such fluctuations are not (e) conclusive evidence: Their predicted shape and magnitude is known to vary with (e&eb) the details of the inflationary potential (whose form is not yet experimentally constrained). |
| (1881) Hence many different observations could be accommodated, and any agreement between theory and observation is at best suggestive. There is also a possibility that other, noninflationary sources might contribute. However, gravitational waves—the propagating ripples in spacetime predicted by Einstein in 1916—might provide (eb) direct evidence of inflation. |
| (1882) During inflation, gravitational fields would also fluctuate because of (e) quantum effects. |
| (1883) These fluctuations would similarly get "frozen in" and re-emerge after inflation as (=) classical gravitational waves, whose amplitude depends only on (=) the energy stored in the field driving (eb) inflation, and not on other model-dependent details [4]. |
| (1884) If the scale of inflation were indeed as high as a few times 10161016 GeV, this background could produce (eb) specific temperature anisotropies in the CMB [5]. |
| (1885) These anisotropies have thus far not been detected by (e) the most recent searches by the WMAP and PLANCK satellites, and their absence was used to (e) put upper limits on the scale of inflation somewhat below 10161016 GeV. |
| (1886) But a more sensitive CMB probe of (e) gravitational waves exists. |
| (1887) Large-wavelength gravitational waves produced by (e) inflation would have alternatively compressed and expanded (e&eb) space, producing (eb) spatial variations in (eb) the intensity of radiation bombarding free electrons in (eb)the early Universe (see, e.g., Ref: [6]). |
| (1888) As photons decoupled from (e) matter and CMB radiation was released, scattering from (e) the undulating sea of electrons would have created (eb) patterns in the polarization of CMB. |
| (1889) Such polarization anisotropies are predicted in (eb) inflationary models to be (=) most visible when (e) looking at the sky on spatial scales comparable to (e&eb) or somewhat larger than the size of (e) the Universe's horizon at (eb) the time when the CMB was created. |
| (1890) This corresponds to (e&eb) about 1 degree of the field of view of the sky. |
| (1891) Over the past decade, various CMB polarization experiments have thus chased the observation of polarization anisotropies on (e&eb) these scales but, so far, also yielded (eb) merely upper limits on a possible gravitational background. |
| (1892) All of this has now changed with (e&eb) the BICEP2 collaboration's [1] report of (e) a signal that has the predicted (eb) characteristics of such a background at precisely the angular scales where (e) the inflationary signal is expected (see Fig. 1, top deleted due to spatial restraints). |

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| (1893) The BICEP2 Collaboration uses (e) a microwave telescope located at the South Pole, with (e&eb) an array of polarization-sensitive detectors pointing at (eb) a small field (a solid angle of about 11 percent of the full sky) with a resolution of approximately 0.50.5 degrees, thus sufficient to probe the relevant angular scales for the detection of (e) polarization anisotropies originating from (e) **inflation**. |
| (1894) BICEP2's search is based on (e) two expected features of primordial gravitational waves. |
| (1895) First, gravitational waves would produce (eb) a characteristic spatial polarization pattern on CMB—a consequence of (e) their quadrupole nature. |
| (1896) One can decompose (e&eb) any spatial polarization pattern into (e&eb) two orthogonal components, a "curl-free" component and (e&eb) a component involving a twisting pattern. |
| (1897) In analogy to electromagnetic modes, these polarization patterns are called "EE modes" and (e&eb) "BB modes," respectively (see Fig. 1). |
| (1898) Scalar fluctuations in the density of (e) matter would produce (eb) EE modes, while (e) gravitational waves would instead generate (eb) a BB-mode pattern. |
| (1899) The second feature is the aforementioned characteristic spatial scale of (e) expected CMB polarization signatures, including (e) wavelength scales comparable to (e&eb) or larger than the horizon scale at the time the CMB was created. |
| (1900) Since inflation would produce (eb) both gravitational waves and (e&eb) scalar density perturbations, BICEP2, like all CMB-polarization experiments, quote their results in terms of (e&eb) the ratio rr of the amplitude of (e) tensor perturbations (associated with (e&eb) gravitational waves) to (e) the amplitude associated with (e&eb) scalar density fluctuations (reflected in (eb) the magnitude of CMB temperature fluctuations measured by (e) experiments like COBE). |
| (1901) A nonzero rr implies (eb) the detection of a BB-mode pattern and primarily measures (eb) the strength of any possible gravitational wave signal which, if (e) it comes from inflation, is then determined by (e) its energy scale (e.g., see Ref. [7] for a review). |
| (1902) BICEP2 reported a nonzero value of r=0.2(+0.07−0.05)r=0.2(+0.07-0.05), on scales ranging from (e) the horizon scale up to 4–5 times larger, with (e&eb) a claimed significance of better than 5σ5σ. |
| (1903) If confirmed, a value of 0.2 corresponds to (e&eb) an energy density during (e&eb) inflation of ∼1.2×1016∼1.2×1016 GeV. |
| (1904) The claimed signal is far larger than expected, partly because (e) the PLANCK satellite, despite not having yet reported its polarization measurements, has (e) previously put an upper limit on (eb) rr (<0.11<0.11). PLANCK's estimate is based on (e) the absence of a gravitational wave signature on (eb) the CMB temperature anisotropy on (eb) its largest scales (PLANCK observes the whole sky and therefore angular scales up to 360360 degrees). |
| (1905) The apparent discrepancy between the PLANCK upper bound and (e&eb) the BICEP2 detection might be a cause for concern. |
| (1906) However, the two bounds probe different spatial scales, and the BICEP collaboration has noted that, while tensor perturbations are expected to be (=) roughly scale independent, some scale dependence of (e) scalar density perturbations could relax (e&eb) the PLANCK constraint and bring (eb) the two results into agreement. |
| (1907) Of greater concern, of course, is the possibility that the signal could be determined by (e) foreground contamination—any source along (e&eb) the line of sight of the detector that could distort or contaminate (e&eb) the primordial CMB signal. |
| (1908) BICEP2 tested for (e) a wide variety of possible foregrounds, claiming (eb) they cannot be responsible for (e) the observed signal. |

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| (1909) One obvious candidate arises from (e) gravitational lensing: The intervening mass of (e) astrophysical bodies could distort (e&eb) the spatial pattern of polarization in a way that that could transform (e&eb) an EE-mode signal into (e&eb) a BB-mode signal. |
| (1910) The BICEP collaboration has analyzed this possibility extensively. Based on (e) measurements of (e) CMB temperature (and hence density perturbations) and on (e&eb) models of the mass distribution, they argue that such effects would only be (=) significant on smaller angular scales than those contributing to their measurements of rr. |
| (1911) But since BICEP2's first announcement on March 17, 2014, a number of scientists have raised another foreground concern due to (e) the possible impact of (e&eb) our Galaxy's dust. |
| (1912) BICEP2 modeled dust contamination scenarios based on (e) a variety of theoretical models, using data from (e) its predecessor BICEP1 and preliminary dust estimates from (e) the PLANCK satellite. |
| (1913) They concluded that dust could not (e) reproduce the magnitude of the observed signal. |
| (1914) However, PLANCK has in the meanwhile released (eb) more complete measurements of variation in polarized dust foregrounds in (eb) many regions of the Galaxy. |
| (1915) Their analysis increases the uncertainty associated with (e&eb) the dust contribution to (e) the BICEP2 result and also accommodates (e) a potentially larger polarized dust contribution. |
| (1916) Some researchers argued that dust polarization at a level 2–3 times higher than that assumed by BICEP2 is possible and could therefore plausibly be responsible for (e) the entire signal (e.g., see Ref: [8]). |
| (1917) While the BICEP2 collaboration continues to maintain that this is extremely unlikely based on (e) their estimates, many other researchers are more skeptical. |
| (1918) In any case, a variety of upcoming measurements, from PLANCK and other detectors are (=) expected to resolve this important ambiguity within the coming year. |
| (1919) Even if the BICEP2 signal does arise from (e) primordial gravitational waves, various consistency checks will be required (e) before it can be definitively argued that the signal arises from (e) inflation. |
| (1920) Inflation is certainly the best-motivated possibility, and the spectral shape of the signal, at least for (e) the five largest-scale modes observed by BICEP2, agrees extremely well with (=) inflationary predictions (see Fig. 1, bottom). |
| (1921) Several results have already appeared that strongly constrain (e) other possible sources, ruling out (e) for example, the possibility of gravitational waves generated by (e) other possible phase transitions after inflation [9] |
| (1922) . This conclusion is reinforced by more careful analyses of the scale dependence of (e) the predicted signal (e.g., Ref. [10]). |
| (1923) But it will nevertheless be important to assure that all possible alternative sources of (e) gravitational waves are ruled out. |
| (1924) If it can be established that the BICEP2 represents (eb) a gravitational wave signal from (e) inflation, there is one additional exciting implication. |
| (1925) Standard calculations suggest (eb) that the gravitational wave signal from inflation is due to (e) quantum fluctuations in the gravitational field, and this appears to require (e) the existence of gravitons—the hypothetical elementary particles that mediate (e&eb) the force of gravity. |
| (1926) Research by Frank Wilczek and myself, based on (e) a dimensional analysis argument [11], suggests that,(eb) independently of the method used to calculate the gravitational wave spectrum, a quantum gravitational origin is (=) required. |
| (1927) If so, the BICEP2 result then implies (eb) that gravity is ultimately a quantum theory, a result of fundamental importance for physics. |

(1928) It is hard to imagine a more important or wider-ranging cosmological observation. If confirmed by upcoming experiments, the BICEP2 results will push the frontiers of empirical physics forward in an unprecedented way, offering (eb) the first direct observation of gravitational waves and of quantum gravity effects as well as a new probe of the earliest moments in the Universe's life and of physics at grand-unified energy scales. This will thus constrain both our understanding of fundamental interactions and the evolution of our Universe as a whole.

**Detection of B-Mode Polarization at Degree Angular Scales by BICEP2 P. A. R. Ade et al. (BICEP2 Collaboration) Phys. Rev. Lett 112, 241101 – Published 19 June 2014**

(1929) Authors report results from the BICEP2 experiment, a cosmic microwave background (CMB) polarimeter specifically designed to search for (e) the signal of inflationary gravitational waves in (eb) the B-mode power spectrum around $\ell \sim 80$.

(1930) The telescope comprised a 26 cm aperture all-cold refracting optical system equipped with (e&eb) a focal plane of 512 antenna coupled transition edge sensor 150 GHz bolometers each with temperature sensitivity of $\approx 300\ \mu K_{CMB}\sqrt{s}$. BICEP2 observed from (e) the South Pole for three seasons from 2010 to 2012.

(1931) A low-foreground region of sky with (e&eb) an effective area of 380 square deg was observed to (e) a depth of 87 nK deg in Stokes Q and U.

(1932) In this paper authors describe the observations, data reduction, maps, simulations, and results. They find an excess of B-mode power over the base lensed-ΛCDM expectation in the range $30<\ell<150$, inconsistent with (e&eb) the null hypothesis at (eb) a significance of $>5\sigma$.

(1933) Through jackknife tests and simulations based on (e) detailed calibration measurements authors show that systematic contamination is (=) much smaller than the observed excess.

(1934) Cross correlating against WMAP 23 GHz maps we find that Galactic synchrotron makes a negligible contribution to the observed signal. We also examine a number of available models of polarized dust emission and find that at their default parameter values they predict power ~(5–10)× smaller than the observed excess signal (with no significant cross-correlation with our maps). However, these models are not sufficiently constrained by external public data to exclude the possibility of dust emission bright enough to explain the entire excess signal. Cross correlating BICEP2 against 100 GHz maps from the BICEP1 experiment, the excess signal is confirmed with $3\sigma$ significance and its spectral index is found to be consistent with that of the CMB, disfavoring dust at $1.7\sigma$. The observed B-mode power spectrum is well fit by a lensed-ΛCDM+tensor theoretical model with tensor-to-scalar ratio $r=0.20+0.07-0.05$, with $r=0$ disfavored at $7.0\sigma$. Accounting for the contribution of foreground, dust will shift this value downward by an amount which will be better constrained with upcoming data sets. Received 4 April 2014 DOI:http://dx.doi.org/10.1103/PhysRevLett.112.241101 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**(1935) Physicist claims to have observed quantum effects of Hawking radiation in the lab for the first time August 16, 2016 by Bob Yirka report**

(1936) This computer-simulated image shows a supermassive black hole at the core of a galaxy. The black region in the center represents the black hole's event horizon, where no light can escape the massive object's gravitational grip. The black …more

(1937) (Phys.org)—Jeff Steinhauer, a physicist at the Israel Institute of Technology, has published a paper in the journal Nature Physics describing experiments in which he attempted to create a virtual black hole in the lab in order to prove that Stephen Hawking's theory of radiation emanating from black holes is correct —though his experiments are based on sound, rather than light. In his paper, he claims to have observed the quantum effects of Hawking radiation in his lab as part of a virtual black hole—which, if proven to be true, will be the first time it has ever been achieved.

(1938) For many years, scientists believed that nothing could ever escape from a black hole. But in 1974, Stephen Hawking published a paper suggesting that something could—particles that are now called Hawking radiation. His idea was that if a particle (and its antimatter mate) appeared spontaneously at the edge of a black hole, one of the pair might be pulled into the black hole while the other escaped, taking some of the energy from the black hole with it—which would explain why black holes grow smaller and eventually disappear. Because such emissions are so feeble, no one has been able to measure Hawking radiation, so researchers have instead tried to build virtual black holes in labs to test the theory. One type of virtual black hole was proposed back in 1981 by Bill Unruh with the University of British Columbia—he suggested that an analogue might be created using water instead of light. He imagined a phonon existing at the edge of a waterfall—as the water speeds up, it begins to move faster than the speed of sound, causing it to be trapped. But if the phonon had an entangled mate that eluded the fall by moving away before getting caught up, it could escape. In this new effort, Steinhauer has built a device based on that idea and in so doing, claims he has observed an analogue of Hawking radiation.

(1939) The experiment consisted of creating an entangled pair of phonons sitting inside a bit of liquid that had been forced (via laser) to move very fast and then observing the action as one of the pair was pulled away as the liquid began to move faster than the speed of sound, while the other escaped—the fluid was a Bose-Einstein condensate of rubidium-87 atoms. After repeating the experiment 4,600 times Steinhauer became convinced that the particles were entangled, a necessity for a Hawking radiation analogue. His findings do not prove Hawking's theory to be true, of course, but they do appear to add a degree of credence that other researchers have thus far not been able to achieve. More information: Jeff Steinhauer. Observation of quantum Hawking radiation and its entanglement in an analogue black hole, Nature Physics (2016). DOI: 10.1038/nphys3863

**(1940) Abstract**

We observe spontaneous Hawking radiation, stimulated by quantum vacuum fluctuations, emanating from an analogue black hole in an atomic Bose–Einstein condensate. Correlations are observed between the Hawking particles outside the black hole and the partner particles inside. These correlations indicate an approximately thermal distribution of Hawking radiation. We find that the high-energy pairs are entangled, while the low-energy pairs are not, within the reasonable assumption that excitations with different frequencies are not correlated. The entanglement verifies the quantum nature of the Hawking radiation. The results are consistent with a driven oscillation experiment and a numerical simulation. Journal reference: Nature Physics search and more info website

Read more at: http://phys.org/news/2016-08-physicist-quantum-effects-hawking-lab.html#jCp

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as

the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t - T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

These equations can be written for

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(13) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.

(14) However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].

(15) Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.

Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

**TIME DELAY**

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period,

which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types system takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t - T_7)F_7(R_7(t - T_7) - N_7(t)M_7((R_7(t) \qquad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7}+\frac{N_7 dG_7}{dR_7}-\frac{dB_7}{dR_7}>\frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7}+\frac{dM_7}{dR_7}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also normalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_7$) and the other derived from the non-delayed or "immediate" processes($I_7$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_7=\frac{N_7 G_7}{2M_7}\frac{d(7nF_7)}{dR_7}(\geq 0) \qquad (7a)$$

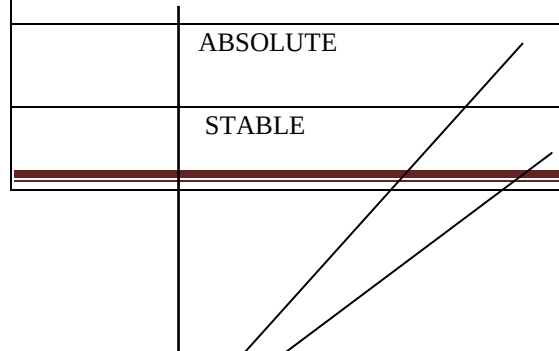
$$I_7=-\frac{N_7 G_7}{2M_7}\frac{d(7nM_7}{dR_7}(\geq 0) \qquad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7>L_7-I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

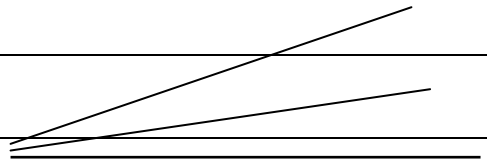
$$P_7>L_7-I_7, \quad \text{when} \quad P_7^2/4>I \qquad (8a)$$

$$P_7>L_7-I_7+\frac{(I_7-P_7^2}{(L_7-I)} \quad \text{when} \quad P_7^2/4<I. \qquad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,





**DELAYED PREDATOR CREDIT TYPES EFFECT WITH ITS WIDE RANGING RAMIFICATION**

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

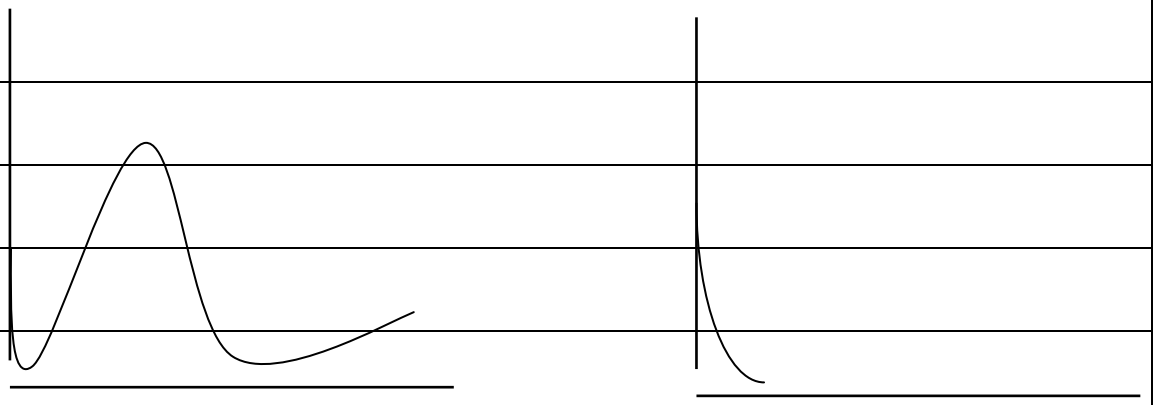
It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay

and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



**TIME DELAY**

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

To scores of problems this relationship is important. Of them some of the examples are as follows:

**New look at the galaxy-shaping power of black holes**

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| (513) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving. |
| (514) "These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science. |
| (515) "Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy." |
| (516) This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return. |
| (517) The remaining matter gets absorbed into the black hole, adding to (e) its mass. |
| (518) The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne. |
| (519) The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution. |
| (520) Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form. |
| (521) "Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara. |
| (522) "It creates (eb) these bubbles, and the bubbles keep the plasma hot. |
| (523) That's what prevents galaxies from (e) becoming even bigger than they are now." |
| (524) Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space. |
| (525) In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene. |
| (526) Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve. |
| (527) "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara. |
| (528) "It's like nature's thermostat, which keeps these galaxies from (e) growing. |
| (529) If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars." |

**A new look at the galaxy-shaping power of black holes**

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| (530) An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the |
| (531) McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said. |

(532) Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.

**Doomed satellite missions**

(533) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.

(534) At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.

(535) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.

(536) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.

(537) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.

(538) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people

(539) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."

(540) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.

(541) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(542) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.

(543) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.

(544) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.

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| (545) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses. |
| (546) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well. |
| (547) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma. |
| (548) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6. |
| (549) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus. |
| (550) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (551) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (552) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (553) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (554) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (555) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (556) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (557) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (558) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (559) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (560) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (561) We know dark matter exists only from (e) its gravitational effects. |
| (562) It holds together galaxies and (e&eb) galaxy clusters. |

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| (563) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (564) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |
| (565) They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (566) X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (567) They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (568) This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (569) An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |
| (570) Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (571) "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (572) The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (573) Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (574) Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (575) "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (576) Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (577) "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New |

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| York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (578) Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories. |
| (579) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen. |
| (580) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer. |
| (581) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants. |
| (582) These were treated as (=) "aware" and "unaware" trials, respectively. |
| (583) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it. |
| (584) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another. |
| (585) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion. |
| (586) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness. |
| (587) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said. |
| (588) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness." |
| (589) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner. |
| (590) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain. |
| (591) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said. |
| (592) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

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| (593) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter. |
| (594) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined. |
| (595) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions. |
| (596) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda. |
| (597) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general. |
| (598) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole. |
| (599) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes. |
| (600) Noda and his colleagues instead studied the change in emissions over (e&eb) time. |
| (601) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection. |
| (602) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1). |
| (603) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole. |
| (604) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved. |
| (605) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp |

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(606) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking “palm tree” organization.

(607) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society

(608) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible?

(609) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem.

(610) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other.

(611) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible.

(612) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges.

(613) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century.

(614) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching.

(615) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom.

(616) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth.

**Spheres on trees**

(617) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration.

(618) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement.

(619) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum.

(620) By visualizing the problem in this way, the researchers noticed that these particular graphs don’t (e) have lots of branches extending from lots of (e) other branches.

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| (621) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure. |
| (622) The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network. |
| (623) As a result, it provides (eb) an important clue for finding the global minimum. |
| (624) It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum. |
| (625) This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large. |

**Thompson's problem 2**

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| (626) By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society |
| (627) The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima. |
| **(628)** The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node. |

**Implications for Thomson's and Smale's problems**

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| (629) Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before. |
| (630) "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around." |
| (631) As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem. |
| (632) "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view." |
| (633) In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges. |
| (634) From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear. |
| (635) "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK. |
| (636) "For larger numbers, there will likely be (=) multiple funnels. |
| (637) Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum." |

(638) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.

(639) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.

(640) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t-T_7)F_7\big(R_7(t-T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t-T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) and the effective prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this

rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Look at the resemblance of String Theory and Abhinava's earlier theory and how the model could be applied: The world of manifestation and differentiated perceptions (Yikalpa) thus extends from one Centre to (eb) the next. **What is meant again is the realisation or the evolution of individual consciousness or general ledger to be (=)the cosmic general ledger is the sole criterion for the Satchidananda state where all dualities cease** Although it is never in fact divorced from (e) the subject who resides there, the ignorant fail (e) to grasp this fact and so, cut off from (e) the Centre, the world of objectivity becomes (=) for them the sphere of Maya. Bhagavatotpala quotes the Light of Consciousness amvitprakdsa): This ever pure experience (suddhanubhava) is variegated by (e) each form [revealed within it] Even so it remains unstained (nirmala) when moving to another. Just as a cloth which is naturally white, once dyed, cannot (e) change colour without [first] becoming white again, similarly the pure power of awareness, (citi) once coloured by (e) form, is pure [again] at the Centre where that form is (=) abandoned and from whence it proceeds to (e&eb) another. In his Essence of Vibration (Spandasarpdoha), K§emaraja explains that the rise and fall of every individual perception in the field of awareness is (=) a specific pulsation of consciousness. This is exactly the individual general ledger transaction lifted to the transaction in cosmic general ledger. During The Initial Instant Of Perception, T Consciousness Is Manifestly Apparent And The Yogi, Participating In Its Plenitude, Observes The Outer World Without Being Attached To Any Particular Divine Body And Sacred Circle Of The Senses Singling It Out From Any Other, Like A Man Who Observes A City From A High Mountain Peak.

He Sees The Outer World Reflected Within His Consciousness Free Of Thought-Constructs And So 'Stamps* The Outer On The Inner While Absorbing The Object And Means Of Knowledge In The Pure Subject Which Grasps Them As The Expansion Of His Own Nature. K§emaraja Says: By Penetrating Into Bhairavimudra, The Yogi Observes The Vast Totality Of Beings Rising From, And Dissolving Into, The Sky [Of Conscious- Ness], Like A Series Of Reflections Appearing And Disappearing Inside A Mirror. Through (e&eb) The Practice Of Bhairavimudra The Yogi Realises That He Is (=) The Substratum Consciousness (Adhisfhatr) Which Both Underlies (e&eb) And Is (=) The Essence Of All Things. He Discovers That Phenomena Have (e) No Independent Existence Apart From (e) Him And So Are, In This Sense, Void. At The Same Time, He Realises That Because (e) All Things Are Consciousness, They Are (=) Far From Unreal. He Views The Outer World Yet Sees (e) It Not. Beyond Both Voidness And Non-Voidness He Penetrates Into (e&eb) The Supreme Abode (Pararn Padam) Of Siva's Consciousness. [The Powers Of The Senses] Endowed With (e&eb) The Attributes Of The Great Union [Between Subject And (e&eb) Object] Whose Form Is The Wakening Of Man's Spiritual Potential (Kuntfalinf), Fill (e&eb) [With Consciousness] The Outer Clatter Of Diversity (Bhedatfambara) Born Of (e) Its Intense Power And Are Then Established In (eb) The Unobscured Abode Of The Void Of Consciousness To Shine [There] Eternally. Thus Residing Beyond (e) Being And Non-Being, The Sole Protector Of The Unity Which Is Tranquil And Expanding [Consciousness], Whose Glory Is All-Embracing And Form Unobscured, Is Called Bhairavimudra. Through (e&eb) The Practice Of Bhairavimudra, The Yogi Unites The Universal Vibration Of T Consciousness With (e&eb) The Individual Pulsation Of Objectivised 'This' Consciousness. The Two Aspects Of Consciousness Are Now In (eb) A State Of Equilibrium Like (=) The Two Pans Of An Evenly Weighed Balance And The Yogi Experiences (e&eb) The Pure Knowledge (Suddhavidya) That: 'I Alone Am (=)All Things'. Thus Becoming (=) The Master Of The Wheel Of Energies He Is (=) Free, Like Siva, To Create And Destroy. When [The Yogi] Is Well Established, Without (e) Wavering, Solely In The Integral Egoity Of His Authentic Nature, The Spanda Principle, And Is

Absorbed In Contemplation (Samavisfa), He Becomes (=) One With It The Doctrine Of Vibration (Tanmaya). Then . . . Dissolving And Creating The Universe By Means Of (er) His Introverted And Extroverted Absorption, He Destroys And Creates (e&eb) All Things Out Of Sarikara, His Innate Nature. [Thus] He Assumes The State Of The Universal Experiencer And Having Absorbed All That Is To Be Experienced From [The Grossest Level] — Earth — To (The Subtlest) — Siva — He Reaches (eb) The State Of The Supreme Subject By (e) Progressively Recognising [His Identity With Him]. Thus, Introverted And Extroverted Absorption Both Lead To (eb) The Recognition Of The Pulsation (Spanda) Of (e) One's Own Consciousness. This Is Where (e) Individual General Ledger Becomes (=) Cosmic General Ledger. At The Level Of Consciousness Corresponding To (e&eb) Siva's Basic State (Saryibhavdvastha), The Alternation From Inner To (e&eb) Outer Is Instantaneously Resolved Into (e&eb) The Vibration Of His Nature. When The Yogi Finally Comes To Be (=) Constantly Aware Of This Reality, His Enlightenment Is (=) Full And Perfect. Freed Of All Means (Anupaya) And Delighting In The Power Of His Bliss (Anandasakti), He Knows And Does Whatever He Pleases. The Yogi Seeking Self-Realisation Must Acquire (e) Mastery Over This Movement. Ksemaraja Stresses That The Doctrine Of Vibration Teaches That Liberation Can Only Be Achieved By (e) First Withdrawing All Sense Activity In Introverted Contemplation (Nimilanasamadhi) To (e) Then Experience The 'Great Expansion' (Mahdvikdsd) Of (e) Consciousness While (e) Recognising This To Be (=) A Spontaneous Process Within (eb) It. This Is Done Through (e&eb) The Practice Of Kramamudra. A Passage From The Now Lost Kramasiitra Explains: Although The Adept's Attention [May Be] Outwardly Directed, (eb) He Enjoys Contemplative Absorption Through (e&eb) The Introverted Aspect Of (e) Kramamudra. Initially He Turns Inward From (e)The Outside World And [Then] From Within [Himself] He Exits Into The Outer World Under The Influence Of His Absorption. There Is No Necessity Of Any Glorification Or Mortification Here. It Is Just Like Joining A Job And Retiring And Realising All That Happened Is Just By His Own Effort And Nothing Else. Thus The Sequence (Krama) In This Attitude (Mudra) [Ranges Through] Both Inner And Outer. The Yogi Must Pervade

The Surface Level Of Awareness (Yyutthana) With The Same Bliss He Experiences Plunging Into The Depths Of Con- Templative Absorption (Samadhi). Note Here That Samadhi Is A Death Like State And Still Being Matter And Information Stored You Experience Consciousness. Submerging Himself And Emerging Repeatedly From Samadhi, He Eventually Recognises That The Unity Of Consciousness Pervades Both States: The Best Of Yogi's, Who Has Achieved A State Of Complete Absorption Even When Risen From Meditation, [Inwardly] Vibrating Like A Drunkard In Blissful Inebriation From The After-Effects Of The Nectar Of Con- Templation, Sees All Things Dissolving In The Sky Of Consciousness Like A Cloud In The Autumn Sky. He Plunges Repeatedly within Himself Divine Body and Sacred Circle of the Senses and Becomes (=) Aware Of His Identity With (e&eb) Consciousness By (e) The Practice Of Introverted Contemplation. Thus Even When He Is Said To Have Risen From (e) Absorption, He Is (=) One With [His] Experience Of It. Ksemaraja Goes On To Explain That This Practice Is Called 'Mudra' Because (e) It Both Fills The Adept With (e&eb) Bliss (Mud) And Is Itself The Bliss Of (e) Consciousness. Moreover, It Dissolves (e) Away (Dra) All Bondage And 'Stamps' The Universe Of Experience With (e&eb)The Seal (Mudra) Of The Fourth State (Turiya) Of Enlightened Consciousness Beyond, And Including, The Three States Of Waking, Dreaming And Deep Sleep. It Is Called 'Krama' Because (e) It Is The Root Source Of (e) All Emanation And All Other Conscious Processes Which Succeed (e&eb) One Another In Ordered Sequence (Krama) And Is, (=)At The Same Time, Their Successive (Krama) Appearance As Well. By The Practice Of Kramamudra The Opposites Fuse And Siva And Sakti Unite. They Yogi Comes To Experience The Simultaneous Pervasion Of All The Lower, Grosser Categories Of Existence By The Higher And The Presence Of The Lower In The Higher. Commencing His Practice In A Low Form Of Bhairavimudra, The Yogi Conjoins The Outer With The Inner; Then, In Kramamudra, He Fills Both The Outer With (e&eb) The Inner And (e&eb) The Inner With (e&eb)The Outer. When He Achieves (eb) Perfection In This Two-Fold Movement, He Attains To (e) The Highest Form Of Bhairavimudra In Which The Two Merge Completely In (eb)The Experience Of (e) The Absolute (Anuttara), Free

Of (e) All Differentiation And (e&eb) Polarities. If He Fails To Maintain Awareness Of This State, He Again Falls Into (eb) Kramamudra Until He Has Finally Completely Merged All The Highest States In (e&eb) The Lower And The Lower In (e&eb) The Higher. He Then No Longer Needs To (e) Resort To (e) Any Means (Anupaya) To (e) Achieve Liberation. All He Says or Does Anything He Perceives or Thinks, Instantly Occasions In (eb) Him The Highest Level Of Consciousness. Thus The Fruit Of Bhairavimudra Is (*=)The Wonder (Camatkara) Or Amazement (Vismaya) That Overcomes (e)The Yogi When (e) He Reaches The Plane Of (e) Union (Yogabhumika) Y Xsl Where (e) All Opposites Merge In (e&eb)The Radiance Of The Great Light Of Consciousness. The Stanzas On Vibration Teach: How Can One Who, As If Astonished, Beholds (e) His Own Nature As (=) That Which Sustains [The Existence Of Everything] Be Subject To (e&eb) This Painful Round Of Transmigration (Kusrti)? The Yogi, Recognising His True Nature To Be (=)The Supreme Subject, Is Astonished To (e) Suddenly Discover That The Individual He Thought He Was, Caught Up (e&eb) In The Trammels Of Thought And Living In A World Enmeshed In (e&eb) The Web Of Time And Space, Does Not Really Exist At All. He Experiences A 'Turning About' (Paravrtti) In (eb)The Deepest Seat Of The Doctrine Of Vibration Consciousness As He Penetrates His True Nature. The Sudden Eruption Of This Intuition (Pratibha) Arouses In Him A Cry Of Amazement As He Transcends All Thought-Constructs And, Perfectly Absorbed In His Own Nature, Is Liberated. The Path To Liberation Essentially, Spanda Doctrine Is Concerned With Two Matters. The First Is To Impart To Those Who Are Fit To Receive The Teachings A Deeper Understanding Of The Ultimate Goal Of Life (Upeya). When We Have Understood What Truly Benefits Us And Is Worth Attaining And What, On The Contrary, Is Of No Real Value But Stands In The Way Of This Attainment, We Can Begin To Make Progress Towards Our Goal. This Is Spanda Doctrine's Second Concern, Namely, To Show The Way In Which We Can Develop Spiritually Through Siva's Grace And The Right Application Of The Means To Realisation That It Teaches. When Both These Aspects Of The Teaching Have Been Correctly Understood And Applied, The Spanda Yogi Achieves A Clear And Permanent

Realisation Of His Goal And Is Liberated, Thus Fulfilling The Ultimate Aim Of The Teaching. The Doctrine Of Vibration Is Not Meant For The Spiritually Dull. It Is Not For The Worldly Whose Consciousness, Clouded By Ignorance, Is As If Dreaming, Even During The Waking State Of Daily Life, The Dream Of Its Own Thought-Constructs. The Teachings Are Meant For Those Who Are Awake (Prabuddha), Those Who, Full Of Faith And Reverence, Are Always Alert And Intent On Discerning The True Nature Of Ultimate Reality. This Reality Is Understood In Three Basic Ways. The First Is Purely Transcendental. The Stanzas Choose This Aspect As The One Which Formally Defines It Most Specifically. Ultimate Reality Transcends All The Opposites, Including Subject And Object. This Does Not Mean, However, That It Is An Unconscious Void, A Mere Absence Of All Existence. In Fact, This Negative Characterization Of Reality (Which Includes Also A Denial Of All That Is Unconscious) Implies A Positive Immanence In Which The Opposites Are United In The Oneness Of Pure Consciousness That Is Equally Siva And Spanda, His Universal Activity. These Two Seemingly Contrasting Aspects Are Reconciled In The Third, Namely, Reality Understood As The Essential Nature Of All Things. Although Universal And Everywhere The Same, It Is Understood To Be The Essential And Specific Nature Of Each Existent As Its 'Own Nature' (Svabhava).

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such

a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

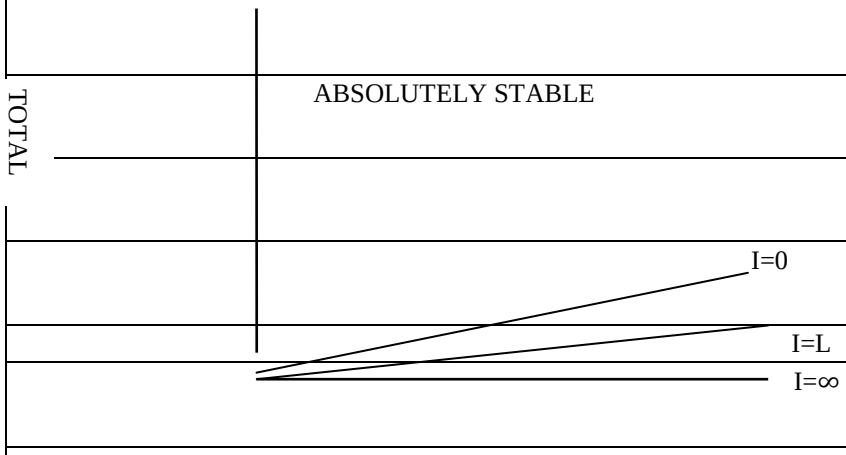
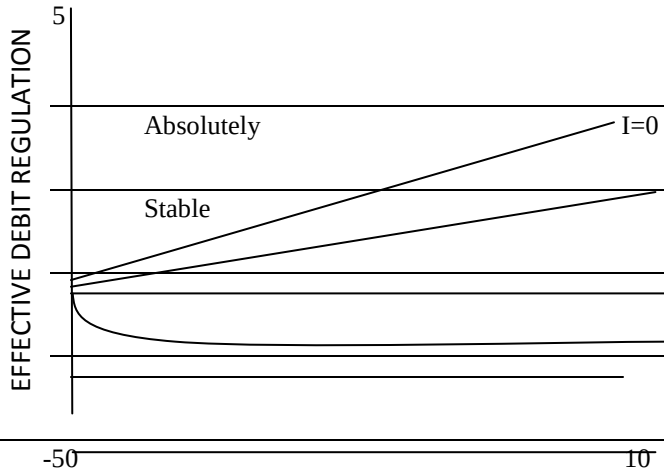



FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

This equation is of paramount importance in the following cases: According to the principles of geometric quantum mechanics, the physical characteristics of a given quantum system can be represented by geometrical features that are preferentially identified in this complex manifold. Here we construct a number of examples of such features as they arise in the state spaces for spin, spin 1, spin and spin 2 systems, and for pairs of spin systems. A study is then undertaken on the geometry of entangled states. A locally invariant measure is assigned to the degree of entanglement of a given state for a general multi-particle system, and the properties of this measure are analysed for the entangled states of a pair of spin particles. With the specification of a quantum Hamiltonian, the resulting Schrödinger trajectories induce an isometry of the Fubini–Study manifold, and hence also an isometry of each of the energy surfaces generated by level values of the expectation of the Hamiltonian. For a generic quantum evolution,

the corresponding Killing trajectory is quasiergodic on a toroidal subspace of the energy surface through the initial state. When a dynamical trajectory is lifted orthogonally to Hilbert space, it induces a geometric phase shift on the wave function. The uncertainty of an observable in a given state is the length of the gradient vector of the level surface of the expectation of the observable in that state; a fact that allows us to calculate higher orders corrections to the Heisenberg relations. A general mixed state is determined by a probability density function on the state space, for which the associated first moment is the density matrix. The advantage of a general state is in its applicability in various attempts to go beyond the standard quantum theory, some of which admit a natural phase-space characterization. **Journal of Geometry and Physics Volume 38, Issue 1, April 2001, Pages 19–53 Geometric quantum mechanics Dorje C. Brody Lane P. Hughston DOI: 10.1016/S0393-0440(00)00052-8** Expansion (unmesa) of [consciousness] or the creative intuition {pratibha) is (=) [experienced] in the interval which divides two [moments] of differentiated perception (vikalpa) It is here that they arise and disappear. The Sastras and Agamas proclaim with reasoned argument that it is (=) free of thought-constructs {Nirvikalpa) and precedes (e) Path to Liberation all mental representations of any object. None can deny that a gap exists between perceptions insofar as two moments of thought are (=) invariably divided. This [gap] is (=) the undifferentiated unity of all the countless manifestations. Similarly, in the outer more objective sphere, where change consists (e) of the alterations in the configurations of manifest appearances (Abhasa\ the transition from one to (e&eb) another corresponds to (e&eb) a phase of pure luminosity that marks (eb) the beginning of one form and the end of another

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^{\infty} \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7 \left( U_7 - V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7 \right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$then given the constraints of $f(T_7)$, $(A1.5) can\ never\ be\ satisfied\ Co$nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7 \lambda_7 e^{-T_7 \lambda_7} + [(U_7 - V_7) M_7 + Y_7] - [((U_7 -$$

V 7M 7−]e−T 7λ 7=0

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$\boldsymbol{Z}_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7) \cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7) \sin(Z_7 T_7) = 0$$

The warrant of Truth is ever elusive when we deconstruct the foundational justifications of our convenient notions about the way the world works. It is easy to confuse what is actually the creation of beliefs with the "discovery of Truth," a common goal of science and theology. Science offers no ontological Scientific Picture of the World but substitutes a number of applied theories as meta-theory. It's close, powerful and useful: it 'works' and allows use to do work. But the still mysterious "true" metatheory may reveal the basis of consciousness. But who knows when we will lift that veil? Permeating the living reality of our culture are certain contagious notions, fads and trends that have the ability to influence the way we think about the nature of Reality and ourselves. Some of them are toxic and can consume you; others are just whims of pop culture infectious and benign as the common cold. **Jung** described the concept of psychic contagion by certain archetypal forces inherent in the human psyche, which manifest in our spiritual lives and belief systems. He spoke of both conscious and unconscious contamination. Notions like this range from simple superstitions to scientific concepts, to urban myths. Notions sweep through our culture and insinuate themselves within its fabric, as fads, whether they are "real" or not, they can be influential. They mood alter us, make us feel we belong. An analysis of these notions is useful in distinguishing a common human phenomenon from any potential "alien influence" which may or may not be exerted on us from an unknown source. **Strange Attractors: Transference, Holography, and an Archetype Burke, J. (2003).**

**Strange Attractors: Transference, Holography, and an Archetype (Doctoral dissertation, Pacifica Graduate Institute, 2003).** In consideration to the Nobel Proze for Chemistry this year(2014), a stability nalysis of the proble in to single molexues would be a contemporaneous question that has to be taken cognisace of and given importance to . Determination of molecular structure by geometry optimization became routine only after efficient methods for calculating the first derivatives of the energy with respect to all atomic coordinates became available. Evaluation of the related second derivatives allows the prediction of vibrational frequencies if harmonic motion is estimated. More importantly, it allows for the characterization of stationary points. The frequencies are related to the eigenvalues of the Hessian matrix, which contains second derivatives. If the eigenvalues are all positive, then the frequencies are all real and the stationary point is a local minimum. If one eigenvalue is negative (i.e., an imaginary frequency), then the stationary point is a transition structure. If more than one eigenvalue is negative, then the stationary point is a more complex one, and is usually of little interest. When one of these is found, it is necessary to move the search away from it if the experimenter is looking solely for local minima and transition structures. The total energy is determined by approximate solutions of the time-dependent Schrödinger equation, usually with no relativistic terms included, and by making use of the Born–Oppenheimer approximation, which allows for the separation of electronic and nuclear motions, thereby simplifying the Schrödinger equation. This leads to the evaluation of the total energy as a sum of the electronic energy at fixed nuclei positions and the repulsion energy of the nuclei. Notable exceptions are certain approaches called direct quantum chemistry, which treat electrons and nuclei on a common footing. Density functional methods and semi-empirical methods are variants on the major theme. **Orlando Alvarez** discusses Polyakov's quantization of the string in the presence of a boundary allowing for an arbitrary topology for the world sheet. In addition to the dynamical conformal factor discovered by Polyakov, there are a finite

number of new degrees of freedom if the surface is more complicated than a sphere or a disc. The quantization of the Liouville theory in an arbitrary topology is also discussed. A one-loop calculation shows that the model is renormalizable if one performs a mass renormalization and an additive field renormalization. The renormalization group equations have a perturbative infrared unstable fixed point in all topologies. Copyright © 1983 Published by Elsevier B.V. **Nuclear Physics B Volume 216, Issue 1, 25 April 1983, Pages 125–184 Theory of strings with boundaries: Fluctuations, topology and quantum geometry Orlando Alvarez DOI: 10.1016/0550-3213(83)90490-X** Quantum geometry (the modern loop quantum gravity involving graphs and spin-networks instead of the loops) provides microscopic degrees of freedom that account for black-hole entropy. However, the procedure for state counting used in the literature contains an error and the number of the relevant horizon states is underestimated. In our paper a correct method of counting is presented by **Marcin Domagala and Jerzy Lewandowski.** Results lead to a revision of the literature of the subject. It turns out that the contribution of spins greater than 1/2 to the entropy is not negligible. Hence, the value of the Barbero–Immirzi parameter involved in the spectra of all the geometric and physical operators in this theory is different than previously derived. Also, the conjectured relation between quantum geometry and the black-hole quasi-normal modes should be understood again. **Marcin Domagala and Jerzy Lewandowski 2004 Class. Quantum Grav 21 5233 doi:10.1088/0264-9381/21/22/014 Black-hole entropy from quantum geometry** Quantum geometry predicts that a universe evolves through an inflationary phase at small volume before exiting gracefully into a standard Friedmann phase. This does not require the introduction of additional matter fields with ad hoc potentials; rather, it occurs because of a quantum gravity modification of the kinetic part of ordinary matter Hamiltonians. Authors draw the cognizance of application of the same mechanism that can explain why the present day cosmological acceleration is so tiny. DOI: http://dx.doi.org/10.1103/PhysRevLett.89.261301

**Inflation from Quantum Geometry Phys. Rev. Lett 89, 261301 – Published 12 December 2002 Martin Bojowald** The loop quantum cosmology of the closed isotropic model is studied with special emphasis on a comparison with traditional results obtained in the Wheeler-DeWitt approach by **Martin Bojowald** .Investigation includes the relation of the dynamical initial conditions to boundary conditions such as the no-boundary or the tunneling proposal and a discussion of inflation from quantum cosmology. **DOI: http://dx.doi.org/10.1103/PhysRevD.67.124023 Loop quantum cosmology, boundary proposals, and inflation Phys. Rev. D 67, 124023 – Published 23 June 2003 Martin Bojowald and Kevin Vandersloot** In loop quantum cosmology, the universe avoids a big bang singularity and undergoes an early and short super-inflation phase. During super-inflation, non-perturbative quantum corrections to the dynamics drive an inflaton field up its potential hill, thus setting the initial conditions for standard inflation. **Shinji Tsujikawa et al** show that this effect can raise the inflaton high enough to achieve sufficient e-foldings in the standard inflation era. They also analyse the cosmological perturbations generated when slow-roll is **violated after super-inflation** and show that loop quantum effects can in principle leave an **indirect signature** on the largest scales in the CMB, with some loss of power and running of the spectral index. **Shinji Tsujikawa et al 2004 Class. Quantum Grav 21 5767 doi:10.1088/0264-9381/21/24/006 Loop quantum gravity effects on inflation and the CMB Quantum geometry from quantum information** is the study of **Florian Girelli and Etera R Livine .**Loop quantum gravity defines the quantum states of space geometry as spin networks and describes their evolution in time. We reformulate spin networks in terms of harmonic oscillators and show how the holographic degrees of freedom of the theory are described as matrix models. This allows us to make a link with non-commutative geometry and to look at the issue of the semi-classical limit of loop quantum gravity from a new perspective. This work is thought of as part of a bigger project of describing quantum geometry in quantum information terms.

**Florian Girelli and Etera R Livine 2005 Class. Quantum Grav 22 3295 doi:10.1088/0264-9381/22/16/011 Reconstructing quantum geometry from quantum information: spin networks as harmonic oscillators.** Quantum geometric mechanics is the topic of study by **Dorje C. Brody Lane P. Hughston** establishing a clear relation between the quantum mechanics and quantum geometry and also the entangled states. The manifold of pure quantum states can be regarded as a complex projective space endowed with the unitary-invariant Fubini–Study metric. Note how differential gheometry concepts with Nunney's model could be entangle so that it is feasibly and viably applied to various type of model including consciousness (See Doctrine of Vibration)

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $\boldsymbol{Z}_7^2$ and $Z_7$in (A2.2):

$W_7^2 (X_7^2$-3 $M_7\ A_7\ X_7$- $M_7^2 (X_7$- $Y_7) + 2\ M_7^3\ A_7$+2 $A_7^2\ M_7^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_7$=1-cos ($Z_7\ T_7$).

In the absence of a delay the system is stable for $A_7 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7 > (X_7 - Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$A_7 > [\ X_7 - Y_7/2\ M_7$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_7, A_7) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_7 \le 2$ *is* of interest First consider the extreme values

$$Q_7(0, A_7) = A_7^2 (X_7 - Y_7)$$

$$Q_7(2, A_7) = (A_7 + 2M_7)^2 (2A_7 M_7 - (X_7 - Y_7))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7 M_7^2 - 3M_7 X_7\right) + W_7\left(3M_7 A_7^2 - 2A_7 X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7 + Y_7\right)$$

$$=2M_7^2 W_7^2 (M_7 + A_7) + M_7 A_7^2 W_7 + (2A_7 M_7 - X_7 + Y_7)(A_7 W_7 + M_7 W_7^2$$

$$+ (2 - W_7)((2M_7 X_7 + M_7 Y_7) W_7 + A_7 (X_7 + Y_7)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $W_7$touches the $W$–axis only at its minimum value in which case (A2.9) shows that

further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$ has no minimum in the range 0 ≤ $W_7$ ≤2.Defining

$Q_7(W_7, A_7) = a_7 W_7^2 + c_7$ *where* $a_7$ $b_7$ *and* $c_7$ Are functions of $A_7$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7 \geq b_7 \geq 0$$

And

$$b_7^2 \geq 4a_7 c_7$$

$$Case(i): Y_7 = 0$$

*Let* $Y_7 = 0$ And try the substitution A= $X_7 = /2M_7$ *in* (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_7 X_7 /$ $2M_7$) is positive for *W<2. Since* $Q_7$ increases with $A_7$(see A2.9), increasing $A_7$will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

*Case* $(ii)$: $Y_7 = y_7 X_7$ ,

$Let\ Y_7 = y_7\ X_7$And try the substitution $A_7 = X_7/(2 + A_7\ y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + n_7 y_7)$ we have

$$W_7^2\ y_7\left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2\left(2 - n_7 + n_7 y_7\right)\right]$$

$$-W_7\ X_7\left[\frac{X_7(1 + n_7 y_7)}{M_7^2(2 + n_7 y_7)^2} + 2\ y_7\left(n_7 - 1\right) + \frac{X_7^2(7 + y_7)}{M_7^2(2 + n_7 y_7)} = 0.\ (A2.12)\right.$$

Since $n_7$ and $A_7$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7), the\ condition (A2.77?) can\ be\ written\ as$

$[\frac{X_7(E_7-7)}{E_7^2}]^2[\frac{X_7^2}{M_7^2} - 4\ y_7\ E_7^2\frac{X_7}{M_7^2} + 4y_7(n_7 - 2)E_7^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$2\ E_7[y_7 E_7 + (y_7\ E_7(7 + y_7)(E_7 - n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - y_7 E_7(7 + y_7)(E_7 - n_7))^{7/2}]$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0 $W_7$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7-7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2-2y_7\frac{(E_7+n_7^2y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right)-4y_7(E_7-n_7)$$

And has the solution

$$\frac{X_7}{M_7^2}>\frac{E_7}{(E_7-7)}[y_7(E_7+n_7^2y_7)+[y_7^2(E_7+n_7^2y_7)^2+4y_7(E_7-7)(E_7-n_7)]^{7/2}]$$

Given $E_7>n_7$this lower limit to $X_7/M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7/M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$in the appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7<n_7$ $(ie.,2A_7M_7^2<X_7-Y_7$)then (A2.7)and (A2.8)show that $Q_7$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2}>2E_7[y_7E_7-(y_7E_7(7+y_7)(E_7-n_7)]^{7/2}]$$

And

$$E_7>n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2 > 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.

If $A_7^2 < 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.+ $(Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$ and $X_7$ and $Y_7$ by $W_7^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7\lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]\, e^{2T_7\lambda_7} + M_7 V_7\, e^{2T_7\lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i}\, Z_7$ *and* separate

(A3.1)into real and imaginary parts

$$Z_7^2 + (M_7 + V_7) Z_7 \sin(Z_7 T_7) - (U_7 M_7 + Y_7) + (M_7(U_7 + V_7) - X_7)$$

$$\mathrm{X} \cos(Z_7 T_7) - 2M_7 V_7 \cos^2(Z_7 T_7) + M_7 V_7 = 0$$

$$(M_7 + U_7) Z_7 - (M_7 + V_7) Z_7 \cos(Z_7 T_7) + (M_7(U_7 + V_7) - X)$$

$$\mathrm{X} \sin((Z_7 T_7) - M_7 V_7 \sin(Z_7 T_7) = 0$$

Using (A.3.3) to eliminate $Z_7$ and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - cos\, Z_7\ T_7$ and $A_7 = U_7 - V_7$ we obtain

$$4\,M_7^2 V_7^2 W_7^4 + 2\,M_7 V_7\ W_7^3[(M_7 - V_7)^2 + A_7(M_7 + V_7) + 2(A_7 M_7 - X_7)]$$

$$+ W_7^2[M_7 - V_7)^2(2\,A_7\ M_7 + Y_7 - X_7) + (2A_7\ M_7 - X_7)^2 - A_7(M_7 - V_7)(2A_7\ M_7 - X_7)$$

$$+ M_7 V_7\,(A_7^2 + 4(X_7 + Y_7))]$$

$$+ W_7[A_7^2 + 2X_7)(2\,A_7\ M_7 - X_7) + 2A_7 Y_7(M_7 + V_7) - M_7 A_7^3] + A_7^2(X_7\ M_7 - Y_7) = 0$$

Being a quartic in $W_7$ (A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7 > 0$$

$$and$$

$$C_7 > (X_7 - Y_7)/2\,M_7$$

*recalling that* $A_7 = U_7$ and defining $C_7 = U_7 + V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0, U_7 V_7 = A_7^2(X_7 + Y_7)$$

$$Q_7\left(2, U_7 V_7\right) = \left(C_7 + 2M_7^2\right)^2 (2C_7 M_7 + Y_7 - X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

## Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

**Module numbered five**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This is the condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference

between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \quad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t - T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\left(R(t)\right)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

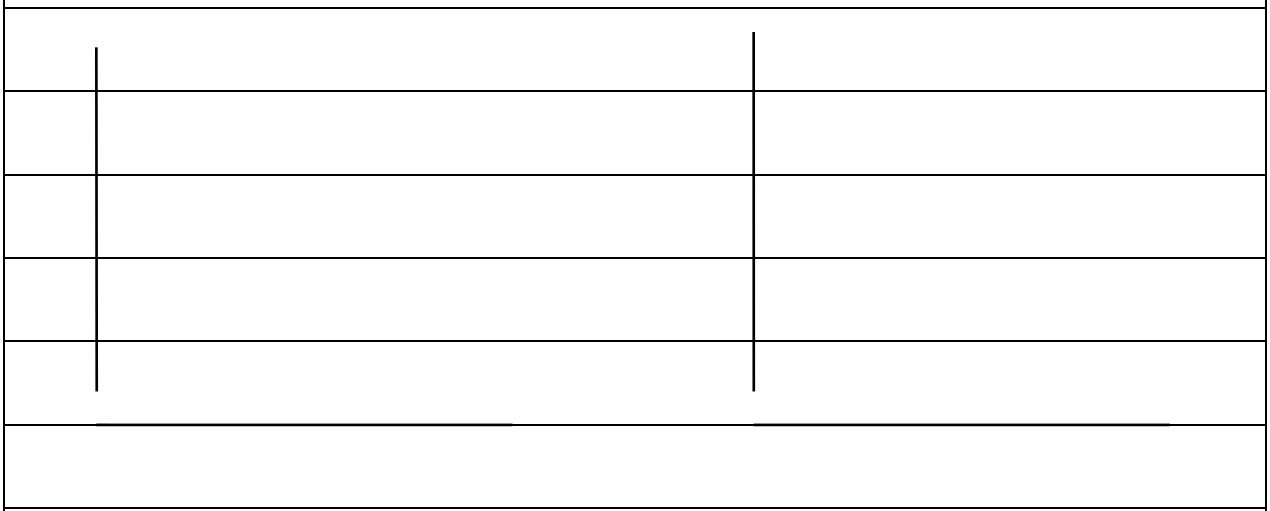
Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time

delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the



TIME DELAY

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period,

which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

As it is usually expressed outside of phenomenological texts, to undergo a conscious experience necessarily means that there is something it is like for the subject to have that experience (Nagel 1974; Searle 1992). This is obviously true of bodily sensations like pain. But it is also the case for perceptual experiences, experiences of desiring, feeling, and thinking. There is something it is like to taste chocolate, and this is different from what it is like to remember what it is like to taste chocolate, or to smell vanilla, to run, to stand still, to feel envious, nervous, depressed or happy, or to entertain an abstract belief. Yet, at the same time, as I live through these differences, there is something experiential that is, in some sense, the same, namely, their distinct first-personal character. All the experiences are characterized by a quality of mineness or for-me-ness, the fact that it is I who am having these experiences. All the experiences are given (at least tacitly) as my experiences, as experiences I am undergoing or living through. All of

this suggests that first-person experience presents me with an immediate and non-observational access to myself, and that consequently (phenomenal) consciousness consequently entails a (minimal) form of self-consciousness. To put it differently, unless a mental process is pre-reflectively self-conscious there will be nothing it is like to undergo the process, and it therefore cannot be a phenomenally conscious process. The mineness in question is not a quality like being scarlet, sour or soft. It doesn't refer to a specific experiential content, to a specific what; nor does it refer to the diachronic or synchronic sum of such content, or to some other relation that might obtain between the contents in question. Rather, it refers to the distinct givenness or the how it feels of experience. It refers to the first-personal presence or character of experience. It refers to the fact that the experiences I am living through are given differently (but not necessarily better) to me than to anybody else. It could consequently be claimed that anybody who denies the for-me-ness of experience simply fails to recognize an essential constitutive aspect of experience. Such a denial would be tantamount to a denial of the first-person perspective. It would entail the view that my own mind is either not given to me at all — I would be mind- or self-blind — or is presented to me in exactly the same way as the minds of others. There are also lines of argumentation in contemporary analytical philosophy of mind that is close to and consistent with the phenomenological conception of pre-reflective self-awareness. Alvin Goldman provides an example: [Consider] the case of thinking about x or attending to x. In the process of thinking about x there is already an implicit awareness that one is thinking about x. There is no need for reflection here, for taking a step back from thinking about x in order to examine it…When we are thinking about x, the mind is focused on x, not on our thinking of x. Nevertheless, the process of thinking about x carries with it a non-reflective self-awareness (Goldman 1970, 96). A similar view has been defended by Owen Flanagan, who not only argues that consciousness involves self-consciousness in the weak sense that there is something it is like for the subject to have the experience, but also speaks of the low-level self-consciousness involved in experiencing my experiences as mine (Flanagan 1992, 194). As Flanagan quite correctly points out, this primary type of self-consciousness should not be confused with the

much stronger notion of self-consciousness that is in play when we are thinking about our own narrative self. The latter form of reflective self-consciousness presupposes both conceptual knowledge and narrative competence. It requires maturation and socialization, and the ability to access and issue reports about the states, traits, dispositions that make one the person one is. To claim that every kind of self-consciousness is conceptual is overly cognitive. Bermúdez (1998), to mention one further philosopher in the analytic tradition, argues that there are a variety of **nonconceptual forms of self-consciousness** that are "logically and ontogenetically more primitive than the higher forms of self-consciousness that are usually the focus of philosophical debate" (1998, 274; also see Poellner 2003). This growing consensus across philosophical studies supports the phenomenological view of pre-reflective self-consciousness. (References: Stanford Encyclopedia, Post Modersnist Theories, Internet Encyclopedia of Philosophy) Korzybski pointed out that there are two ways to slice easily through life; to believe everything or to doubt everything. Both ways save us from thinking. So, from time to time you have to give yourself a Virtual Reality Check. The hallmark of postmodern philosophy has been disbelief or skepticism of all "metanarratives," or translations of reality. Postmodernism has even turned its profound skepticism on such important humanist concepts as "objective truth" and reason.Yet, for a deconstructionist postmodern society, individually we are still riddled with superstition and gullibility, and open to manipulation through our belief systems as any politician, philosopher, clergy, or salesperson will attest. Further, most people are painfully naive when it comes to even the simplest scientific understanding. Most of us don't have a clue about the fundamental nature of physical reality or our own psychological nature, and our ability to be fooled by our senses and mind, or methods of social and technological persuasion. We may think we know, but the penetration is deeper than we dare imagine. Nature rejects the naiveté that seeks absolute truth. We are beginning to realize, individually and culturally, that "realities" are all human constructions. The task becomes one of "catching ourselves in the act" of creating our own "reality" from the flow of events. Human truth is always an engagement of mind with experience.. We don't need to fear the collapse of our

personalistic belief system (the "box" we live in), nor our belief in absolute truth. **A strong desire to engage in the "quest for uncertainty" complements our anxiety that perhaps there is no absolute, objective ground to reality**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t-T_8)F_8(R_8(t-T_8) - N_8(t)M_8((R_8(t) \qquad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8} + \frac{N_8 dG_8}{dR_8} - \frac{dB_8}{dR_8} > \frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8} + \frac{dM_8}{dR_8}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_8 = \frac{N_8 G_8}{2M_8}\frac{d(8nF_8)}{dR_8} (\geq 0) \quad (8a)$$

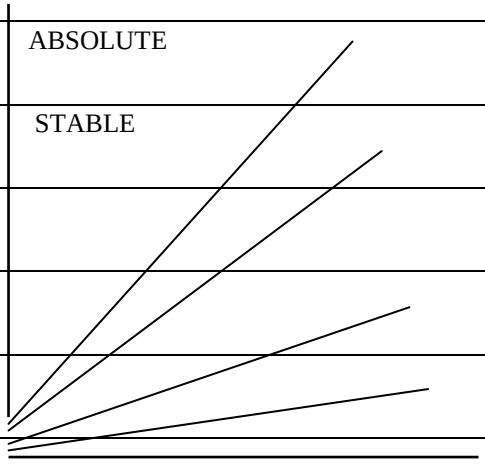
$$I_8 = -\frac{N_8 G_8}{2M_8}\frac{d(8nM_8}{dR_8} (\geq 0) \quad (8b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8 > L_8 - I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_8 > L_8 - I_8$, when $P_8^2/4 > I$ (8a)

$P_8 > L_8 - I_8 + \frac{(I_8 - P_8^2}{(L_8 - I)}$ when $P_8^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

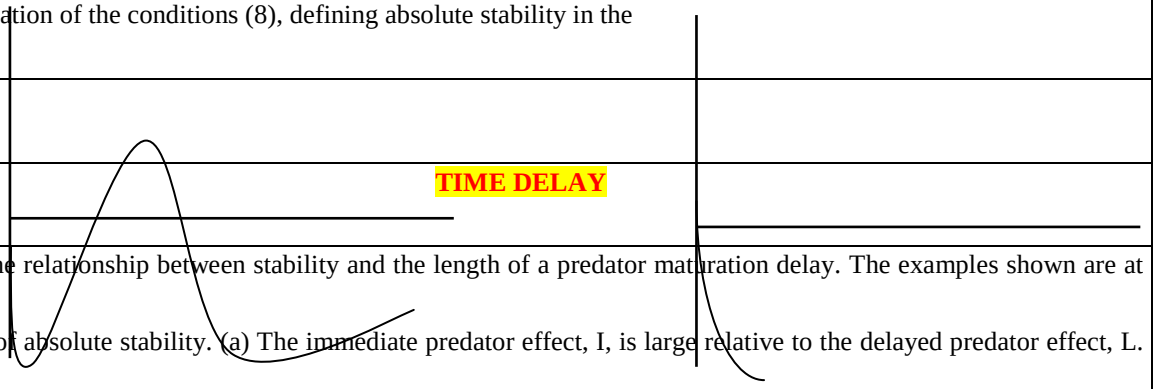



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this

rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior , we must note that the model could be applied to psychological systems too, albeit we are still not sure of the fact how the wave function collapses thereby due to perception.

Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility. He privatizes the sublime, having learned from the fate of his predecessors that the public can never be more than beautiful."(11) Notice is made here of what is indeed the telling feature. Levinas is praised, for example, for his alarm when "the social will [is] sought in an ideal of fusion . . . the subject losing himself in a collective representation, in a common ideal. . . . It is the collectivity which says 'us,' and which, turned toward the intelligible sun, toward the truth, experience, the other at his side and not face to face with him. The tellingly titled D'un ton apocalyptique contains passages that are particularly excited: These people situate themselves outside the ordinary, but they have in common this: they describe themselves as having an immediate and intuitive relation with mystery. And they want to attract, to seduce, and lead others to the mystery, through mystery. This agogic function of the leader of men, of the duce, of the Führer, of the leader, places him above the crowd that he manipulates with the aid of a small numbers of adepts joined together in a sect with a secret language, a clique or a small party with its

ritualized practices. The mystifiers pretend to have exclusive access to the privilege of a secret mystery. . . . The revelation or the unveiling of the secret is something that they jealously reserve for themselves. Jealousy is here a major characteristic. There is not narcissism and non-narcissism; there are narcissisms that are more or less comprehensive, generous, open, extended. What is called non-narcissism is in general but the economy of a much more welcoming, hospitable narcissism, one that is much more open to the experience of the other as other. I believe that without a movement of narcissistic reappropriation, the relation to the other would be absolutely destroyed; it would be destroyed in advance. The relation to the other--even if it remains asymmetrical, open, without reapporpriation--must trace a movement of reappropriation in the image of oneself for love to be possible, for example. Love is narcissistic. Beyond that, there are little narcissisms, there are big narcissisms. . . .(23) **Justice of the Pieces Deconstruction as a Social Psychology Douglas Collins** Sabine Wilke cogently argues for the rise of the aesthetic dimension in philosophical writing by showing how the metaphorical and poetic style of Heidegger, the principles of phrase configurations with Adorno, or the topographic writing of derrida replace through their very form something what traditionally would have been subject to discursive exposure The concern shifts from telling to showing, showing first and foremost that something remains untold. An exemplary book in this aesthetic drive is Avital Ronell's Telephone Book where the playful layout and explicit break with traditional design interrupt the sequential logic of the paragraphs thereby performing the very argument at stake: open up new circuits of meaning and signification. The engagement with form and publishing formats is only consequential since media theory as exposed by Walter J. Ong, Eric A. Havelock, Parry Miman or Marshall McLuhan besides many others made it clear that different communication technologies condition the mode of thought of a given culture The fractals of Benoit Mandelbrot or the butterfly shaped Lorenz attractor are the popular icons of this turn in the sciences. For communication studies and philosophy, Vilém Flusser was one of the pioneer thinkers reflecting the rise of what he called “techno-images.” The photograph is the paradigmatic case which condenses the whole crises of Western

thought and culture on its grainy surface. It subverts the linear logic of writing.6 Flusser's central diagnosis detects a shift from the grand narratives associated with the literate mind and Modernity towards creating and envisioning the world through image based practices. Writing himself in a succinct, visually inspired style, Flusser contents that the web of meaning, spun throughout the centuries by our textual practices, has become transparent. We discover the void between the letters through which the world escapes: "We are alienated from the world of text, precisely because we see through them as our own product."7 As a consequence we are confronted with a world of dissociated facts and information bits. Flusser hails this fragmented world as the advent of the zero dimensional point-universe that will provide the new raw material for our image based practices. Techno-images are nothing else than a conscious, inherently artistic drawing-together of point-elements as exemplified first and foremost by the silver-grains of the photograph but also by more recent pixel based techno-images. Similar to blueprints, designs or statistical curves, the weather map for example visualizes and contracts a tremendous amount of meteorological data from widely distributed stations into a single image of high- and low pressure

areas, temperatures, wind speeds, etc. Techno-images "compute" in this sense the universe of unimaginable, abstract information-bits into a concrete surface – the image. Echoing Paul Feyerabend, these image practices have to be understood as creative practices where the truth of scientific models has been replaced by its aesthetic quality **Aesthetics, Writing, Networked Computers Jörg Muller Dissertation Submitted to the Division of Media and Communications of The European Graduate School in Candidacy for the Degree of Doctor of Philosophy**

**Lothar Schäfer (1) and Sisir Roy (2) in** Quantum Reality, the Importance of Consciousness in the Universe, the Discovery of a Non-Empirical Realm of Physical Reality, and the Convergence with Ancient Traditions of Indian and Western Philosophy investigate metaphysical aspects that include 1) the discovery of a non-empirical part of physical reality in a realm of potentiality; 2) the emanation of the empirical world out of a realm of non-material forms; 3) the discovery that the nature of physical reality is that of an indivisible Wholeness – the One; and 4)

panpsychism: the possibility that the One is aware of its processes like a Cosmic Consciousness. The convergence of powerful traditions of seemingly disparate cultures is particularly important to point out in the present process of globalization, when a unifying view is needed to avoid controversy and conflict. To quote them in extenso : This is the nonlocality of the quantum world. Menas Kafatos and Robert Nadeau (1990) have drawn a remarkable conclusion from this phenomenon: If reality is nonlocal, the nature of the universe is that of an undivided wholeness. Because our consciousness has emerged from this wholeness and is part of it, it is possible to conclude that an element of consciousness is active in the universe: a Cosmic Consciousness. Pre-reflective self-consciousness is pre-reflective in the sense that (1) it is an awareness we have before we do any reflecting on our experience; (2) it is an implicit and first-order awareness rather than an explicit or higher-order form of self-consciousness. Indeed, an explicit reflective self-consciousness is possible only because there is a pre-reflective self-awareness that is an on-going and more primary self-consciousness. Although phenomenologists do not always agree on important questions about method, focus, or even whether there is an ego or self, they are in close to unanimous agreement about the idea that the experiential dimension always involves such an implicit pre-reflective self-awareness. In line with Edmund Husserl (1959, 189, 412), who maintains that consciousness always involves a self-appearance (Für-sich-selbst-erscheinens), and in agreement with Michel Henry (1963, 1965), who notes that experience is always self-manifesting, and with Maurice Merleau-Ponty who states that consciousness is always given to itself and that the word ‘consciousness’ has no meaning independently of this self-givenness (Merleau-Ponty 1945, 488), Jean-Paul Sartre writes that pre-reflective self-consciousness is not simply a quality added to the experience, an accessory; rather, it constitutes the very mode of being of the experience: This self-consciousness we ought to consider not as a new consciousness, but as the only mode of existence which is possible for a consciousness of something (Sartre 1943, 20 [1956, liv]). The notion of pre-reflective self-awareness is related to the idea that

experiences have a subjective 'feel' to them, a certain (phenomenal) quality of 'what it is like' or what it 'feels' like to have them.

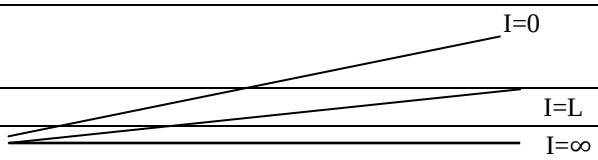
Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

TOTAL

ABSOLUTELY STABLE

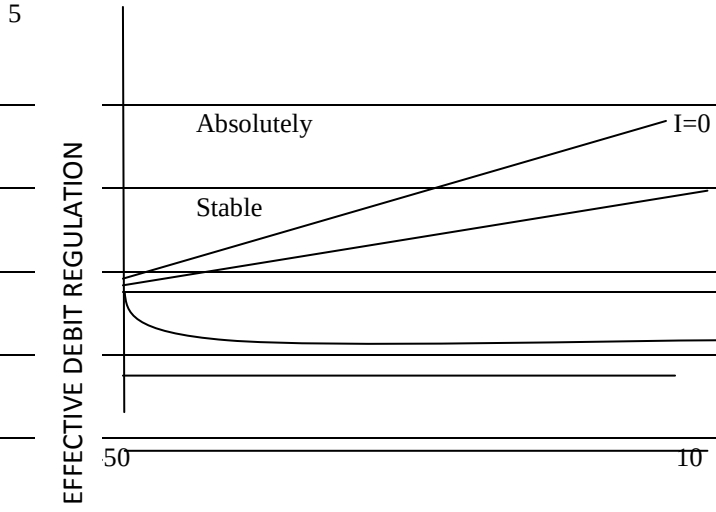



**DELAYED DEBIT EFFECT**

FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

## Utility Analysis And neural Network

Caroline Williams investigated affective density of the political and its effect on our understanding of political subjectivity. Taking up Spinoza's challenge to think about affect beyond corporeal embodiment, he argues that there is a modality of affectivity that cannot simply be inscribed within the borders of subjectivity. theorising affect as an impersonal force anchored in a relational ontology that gives due recognition to the circulation of affects, as well as to their ambivalent structure in creating sites of identification, and he utilises this ontology to reflect on the dynamic of the political and the shape of political subjectivity. Williams argues that Spinoza's philosophy (through ideas of conatus and imagination) offers the conceptual resources to reconfigure the composition of affective subjectivity as a transindividual social bond and as an unconscious dynamic of ethico-political existence. **Subjectivity (2010) 3, 245–262. doi:10.1057/sub.2010.15 Affective processes without a subject: Rethinking the relation between subjectivity and affect with Spinoza Caroline Williams** Phenomenological Time bears ample evidence and infallible observatory and apodictic evidence albeit with a sense of chagrined circumvention and defeated disgruntlment circularity of both psychological and phenomenological interpretations of art which are sometimes insightful, but most often work on the logic of cyclical time and unavoidable drives and representations played out as a ‘natural’ order: where ‘men can’t avoid being boys’ a circularity which valorizes a helplessness/truthfulness conundrum. One begins with a classical phenomenological fallacy with Duchamp’s Fountain, does the urinal keep returning or do we keep returning to the urinal? Put in this way, it is a male question, something that phenomenology conveniently conceals when positing the “I”, or indeed, the “we”. But we must be careful not to simply assume that the phenomenological first person is white, heterosexual, able-bodied and male, providing a foundation upon which knowledge should be based as a universal condition. Deleuze’s second critique of time is linked to Kant. This model of time extends not only to phenomenology but to the phenomenological interpretation

of art and is seen as a straight line. In the Critique of Pure Reason, Kant cuts up circular time and reassembles it as a series of sensory experiences which memory processes and which constitute subjectivity. Whereas the circular model of time (and interpretation) inscribes fatalism into the subject and the world, the phenomenological method is in danger of solipsism, where the subject constitutes the world. Much of embodied philosophy and enactivist theories of cognition go down this line with a plethora of art historical analyses by the likes of Rosalind Krauss, Michael Fried and others taking the lead from Merleau-Ponty. These approaches successfully challenge the **homunculus model of a Cartesian mind-body** dualism where there is a supervisor in the mind who processes experience and even consciousness and to whom all representations are directed. But phenomenology is mainly concerned with perceptual consciousness, the processing of colors, shapes, forms, motion in art and the world. The senses and sensorimotor processes are the fundamental pathways by which the interior is connected to the exterior in a chiasm of self-constitution in the world. The problem with this is that we explore the art through our sense perceptions, and what does art reveal to us?—our sense perceptions! Such solipsistic approaches have real problems relating complex and abstract conceptual production to bodily and sensorimotor contingencies, and they nearly always posit the subject as the site and object of discovery in the world. To be fair, phenomenology has moved on to rather more sophisticated models of intersubjectivity, taking into account the specificities of different kinds of bodies and ways of being beyond the white, male heterosexual as the common denominator, yet many naïve approaches to art perpetuate the notion that perceptual, rather than conceptual processes and evolutionary neural patterns are primary, fundamental ways to understand the art experience In fact, these aspects are only part of art production and reception, aspects which are often filtered out or made known to us in order for us to parody such trigger responses, as in much of Duchamp's work, which holds in contempt such retinal approaches to art. If we take the **subject of abstract expressionism** in the 1950s or later minimalist art, art history, teaching and interpretation still insists on putting the act of perceiving as fundamental to the constitution not only of the phenomenological first

person but also of art itself. The male act of usually standing up, which consists of directing a stream with the hand(s) into a container or on a flat surface, sand or snow in order to leave patterns and traces (the usual child's prank) underlies the logic of macho 1950s Abstract Expressionist culture, especially that of Jackson Pollock, so effectively parodied by Carolee Schneeman's Vagina Scroll, 1975. Here, painting or writing is a kind of a discharge which invests the **object and space with** the phenomenological "I", disciplined by existential struggle and reflected back by the traces, signatures and gestures left behind in the work of art, rediscoverable as an element in the making of other subjectivities.It is precisely this kind of solipsism which is blind to the art event as something well in excess of re-establishing habits of subjective self-identity through the production and interpretation of art. The attempt to reinforce a phenomenological identity upon the event by personalizing sense experience is unable to see art as something which remodels the phenomenological 'I' and neural plasticity itself. We learn new technologies and redraw our neural pathways through such new and spontaneous experiences which draw us into intersubjective worlds. Duchamp's Fountain and Deleuze's Repetition and Différence Grégory Minis sale (I have made some changes) The relation of different groups to deconstruction is settled by its relation to groups. An American enthusiast--one who has had a role in the development of our affection for Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^{\infty} \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$then given the constraints of $f(T_8)$, $(A1.5) can\ never\ be\ satisfied\ Co$nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8 \lambda_8 e^{-T_8 \lambda_8} + [(U_8 - V_8) M_8 + Y_8] - [((U_8 - V_8) M_8 -] e^{-T_8 \lambda_8} = 0$$

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$\boldsymbol{Z}_8^2 + M_8 Z_8 \sin(Z_8 T_8) - (A_8 M_8 + Y_8) + (A_8 M_8 + X_8)\cos(Z_8 T_8) = 0$$

$$(A_8 M_8) Z_8 - M_8 Z_8 \cos(Z_8 T_8) + (A_8 M_8 - X_8)\sin(Z_8 T_8) = 0$$

Where $A_8 = U_8 - V_8$Using (A2.3), substitute for $\boldsymbol{Z}_8^2$ and $Z_8$in (A2.2):

$W_8^2 (X_8^2 - 3 M_8 A_8 X_8 - M_8^2 (X_8 - Y_8) + 2 M_8^3 A_8 + 2 A_8^2 M_8^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_8 = 1 - \cos(Z_8 T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8 / 2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_8 \le 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2\ (\mathrm{X}_8 - Y_8)$$

$$Q_8(2, A_8) = \left(A_8 + 2\,\mathrm{M}_8\right)^2 (2A_8\,\mathrm{M}_8 - (\mathrm{X}_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2\left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8\left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8\left(2X_8 + Y_8\right)\right)$$

$$+2A_8\left(X_8 + Y_8\right)$$

$$=2\,M_8^2 W_8^2 (M_8 + \mathrm{A}_8) + M_8\, A_8^2 W_8 + \left(2A_8 M_8 - X_8 + Y_8\right)(A_8 W_8 + M_8 W_8^2$$

$$+ (2- W_8)((2\,M_8\,X_8 + M_8\,Y_8)\,W_8 + A_8\,(X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $\mathrm{W}_8$touches the $W$ –axis only at its minimum value in which case (A2.9) shows that further increases in *A* will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range $0 \le W_8 \le 2$.Defining

$Q_8\left(W_8, A_8\right) = a_8 W_8^2 + c_8$ *where* $a_8\, b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \boxed{2} - b_8 \boxed{0}$$

And

$$b_8^2 \boxed{4} a_8 c_8$$

$$Case(i): Y_8 = 0$$

$Let\ Y_8 = 0$ And try the substitution A= $X_8 = /2M_8\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8 /$ $2M_8$) is positive for *W<2. Since* $Q_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\ X_8$ ,

$Let\ Y_8 = \ y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 + \ n_8 y_8)$ we have

$$W_8^2\, y_8 \left[X_8 n_8 \frac{(1 + n_8 y_8)}{(2 + n_8 y_8)} + M_8^2 \left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8 X_8 \left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2}+2y_8\left(n_8-1\right)+\frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)}\right]=0.\ (A2.12)$$

Since $n_8$ and $A_8$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8=(2+n_8 y_8)$, *the condition*$(A2.88?)$*can be written as*

$[\frac{X_8(E_8-8)}{E_8^2}]^2[\frac{X_8^2}{M_8^2}-4y_8 E_8^2\frac{X_8}{M_8^2}+4y_8(n_8-2)E_8^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8/W_8^2)$ can be expressed as

$$2E_8[y_8E_8+(y_8E_8(8+y_8)(E_8-n_8))^{8/2}]>\frac{X_8}{M_8^2}>2E_8[y_8E_8-y_8E_8(8+y_8)(E_8-n_8))^{8/2}]$$

Thoughts on the Social and Political Implications of Correlationism .Lion Mirror Truth be told, as my thought has evolved the issue of correlationism had fallen off the radar for me. Somehow the debate had come to seem too "philosophical" to me, too "scholastic", too remote from what interests me: understanding why social assemblages are organized as they are, how power functions in social assemblages, and what we might do to address that power and change things. Somehow the question of whether or not we can get out of the correlation between thinking and being just came to seem remote from these sorts of issues. Somehow it seemed too epistemological. ASIDE: Numerous discussions over the years have led me to believe that the debate over correlationism is poorly understood (or maybe I just don't understand it). On countless occasions I've heard people say "of course we must relate to things in order to know them." Well yeah, of course! I don't think this is what the critic of correlationism is getting at. It seems to me that correlationism is something more robust than the theses that we must relate to something to know it.

Correlationism instead seems to require the theses that thought and being are indiscernible. Put more concretely, the correlationist is someone who argues that we either a) can never tell whether being is merely a construction of our thought (weak correlationism), or b) who argues that thought actually constructs being (strong correlationism). In other words, correlationism is another name for idealism. One can hold that we must relate to something in order to know it without being a correlationist. As an aside I should also add that I am a correlationist about some things. For example, I think money is something constructed by society and is therefore a strong correlationist when it comes to money. At any rate, for a long time I'd become rather indifferent to debates about correlationism and philosophies of access. I had learned the lessons of speculative realism– which I could have also learned, I think, from Deleuze and Guattari, the new materialist feminists, actor-network theorists such as Latour –and had moved on. However, occasionally you come across a tone of phrase that pitches something in a different light. In The Cut of the Real, Katerina Kolozova writes, …the political problem of contemporary philosophy identified by the 'new realists' is, in fact, the product of a more fundamental epistemic problem. In his book After Finitude, Quentin Meillassoux calls this problem 'correlationism' and identifies it as an essentially post-Kantian legacy, which continues to dominate and limit philosophy. As a matter of fact, correlationism lies at (eb) the heart of postmodern theory and consists in the premise that thought can only 'think itself,' that the real is inaccessible to knowledge and human subjectivity, and that there is nothing but discursive constructs that fully determine thinking and that are meth0dologically accounted for all the way down. (1 – 2) Thought thinking only itself: Thought only encountering itself. In the jargon of postmodern and poststructuralist lingo, this would be the thesis of infinite semiosis, where signs ("thoughts") only ever relate to other signs. Within this framework, discursivity comes to be the hegemonic framework defining all of being. At the level of politics and social theory more generally, if the correlationist thesis is true the

consequences are clear: all social phenomena are discursive and all solutions to social and political problems will be discursive. The sole sphere of the political will be the discursive and all questions of politics will be questions of speech-acts and interpretation. The problem here is not that many theorists recognize that the discursive and semiotic plays an important role in the social and the political. It does and I've repeated this tirelessly. The problem is with what happens when thought or the semiotic becomes a hegemon, an "all", foreclosing our ability to recognize other forms of power. What I've wanted to say is that not all power functions discursively. In my last post and elsewhere I spoke of some other forms of politics: Thermopolitics: The politics surrounding **energy in the form of calories and fuels such as gasoline** and coal, and how our life and our very bodies are structured by energy dependencies and by being trapped in particular distributive networks that render these forms of energy available. I'm being quite literal when I speak of energy, talking about the effects, for example, of the absence of food in certain educational environments on cognition, for example; and am generally hostile to metaphorical extensions of the concept of energy which I see as erasing the dimension of real materiality. Geopolitics: The role that features of natural and built geography such as mountain ranges, rivers, oceans, soil conditions, roads, housing design, etc., play in the form that social relations take and how they impact individual bodies. Chronopolitics: The way in which the structurations of time organize what is possible for us. For example, the structuration of the working day, how much we can say and comprehend at any given time, the impact of things like the invention of the clock, etc. **Oikopolitics: This would be the domain of political economy described so well by Marxists**. So five different types of politics: Semiopolitics (or what currently dominates critical theory), thermopolitics, geopolitics, chronopolitics, and oikopolitics. No doubt there are other sites of the political or political struggle that we could speak of, but this is a good start. Also, it should be obvious that these aren't exclusive domains, but are entangled in all sorts of important

ways. For example, something might take place at the level of semiopolitics (speech, law, rhetoric, norms, and communication) that has all sorts of impact at the level of thermopolitics. Congress might decide to cut programs that fund school meal programs. This, in turn, will have a thermodynamic impact on those students that go without the calories they need developmentally and cognitively to function in a particular way. There is an entanglement here of semiopolitical and thermopolitical domains. The young student here has been constrained both at the level of **semiotic phenomena and thermodynamic structures**. The point is that if true, semiotic intervention (speech-acts, protests, interpretations, deconstructions, etc) will not be an appropriate response to all political problems because social formations are not entirely structured by the semiotic. The child in that school does not suffer from a lack of the right signs, but from a lack of calories needed to run the engine of his thought and body. Certainly semiotic interventions might be needed to render that energy available, but it is the energy itself that is at issue and the absence of that energy that forms the spider web entangling him in his position. A correlationist perspective tends to erase this as even being a site of the political. **Larval Subjects January 28, 2014 Ontological Anarché Posted by larval subjects under Uncategorized January 27, 2014 Roden on Pluralism Over at enemy industry Roden has an excellent post up on the pluralism discussion. Check it out here. January 25, 2014 Different Senses of Pluralism and Ontology Posted by larval subjects under Uncategorized Responding to one of my comments over at his blog, Jeremy Trombley:** I have deleted restricted, reformulated some sentences .Kindly bear with me. Intention is to make the discursive enucleation consistent with the subject in question**.** Recourse to epistemology of the quantum mechanics generic and fervor of string theory is a sine qua non for the further comprehension of the theory and ontology .How would the world appear to us if its ontology was that of classical mechanics but every agent faced a restriction on how much they could come to know about the classical state? We show that in most respects

it would appear to us as quantum. The statistical theory of classical mechanics, which specifies how probability distributions over **phase space evolve** under Hamiltonian evolution and under measurements, is typically called Liouville mechanics, so the theory we explore here is Liouville mechanics with an epistemic restriction. The particular epistemic restriction we posit as our foundational postulate specifies two constraints. The first constraint is a classical analog of Heisenberg's uncertainty principle; the second-order moments of position and momentum defined by the phase-space distribution that characterizes an agent's knowledge are required to satisfy the same constraints as are satisfied by the moments of position and momentum observables for a quantum state. The second constraint is that the distribution should have maximal entropy for the given moments. Starting from this postulate, authors derive the allowed preparations, measurements, and transformations and demonstrate that they are isomorphic to those allowed in Gaussian quantum mechanics and generate the same experimental statistics. They argue that this reconstruction of Gaussian quantum mechanics constitutes additional evidence in favor of a research program wherein quantum states are interpreted as states of incomplete knowledge and that the phenomena that do not arise in Gaussian quantum mechanics provide the best clues for how one might reconstruct the full quantum theory. I: http://dx.doi.org/10.1103/PhysRevA.86.012103 **Reconstruction of Gaussian quantum mechanics from Liouville mechanics with an epistemic restriction Phys. Rev. A 86, 012103 – Published 10 July 2012 Stephen D. Bartlett, Terry Rudolph, and Robert W. Spekkens.** Deleuze in fact characterizes modes of existence, with their powers and capacities. The answer is this: Deleuze approaches modes of existence, ethically speaking, not in terms of their will, or their conscious decision making power (as in Kant), nor in terms of their interests (as in Marx, for example), but rather in terms of their drives. For Deleuze, conscious will and preconscious interest are both subsequent to our unconscious drives, and it is at the level of the drives that we have to aim our ethical analysis. Here, I would like to

focus on two sets of texts on the drives taken, not from Nietzsche and Spinoza, but rather from Nietzsche and Leibniz (Leibniz being one of the first philosophers in the history of philosophy to have developed a theory of the unconscious). The first set of texts comes from Nietzsche's great early book entitled Daybreak, published in July 1881. Nietzsche first approaches the question of the drives by giving us an everyday scenario: "Suppose we were in the market place one day," he writes, "and we noticed someone laughing at us as we went by: this event will signify this or that to us according to whether this or that drive happens at that moment to be at its height in us—and it will be a quite different event according to the kind of person we are. One person will absorb it like a drop of rain, another will shake it from him like an insect, another will try to pick a quarrel, another will examine his clothing to see if there is anything about it that might give rise to laughter, another will be led to reflect on the nature of laughter as such, another will be glad to have involuntarily augmented the amount of cheerfulness and sunshine in the world—and in each case, a drive has gratified itself, whether it be the drive to annoyance, or to combativeness or to reflection or to benevolence. This drive seized the event as its prey. Why precisely this one? Because, thirsty and hungry, it was lying in wait" (D 119). This is the source of Nietzsche's doctrine of perspectivism ("there are no facts, only interpretations"), but what is often overlooked is that, for Nietzsche, it is our drives that interpret the world, that are perspectival—and not our egos, not our conscious opinions. It is not so much that I have a different perspective on the world than you; it is rather that each of us has multiple perspectives on the world because of the multiplicity of our drives—drives that are often contradictory among themselves. "Within ourselves," Nietzsche writes, "we can be egoistic or altruistic, hard-hearted, magnanimous, just, lenient, insincere, can cause pain or give pleasure" (Parkes, pp. 291-292). We all contain such "a vast confusion of contradictory drives" (WP 259) that we are, as Nietzsche liked to say, multiplicities, and not unities. Moreover, these drives are in a constant struggle or combat with each other:

my drive to smoke and get my nicotine rush is in combat with (but also coexistent with) my drive to quit. This is where Nietzsche first developed his concept of the will to power—at the level of the drives. "Every drive is a kind of lust to rule," he writes, "each one has its perspective that it would like to compel all the other drives to accept as a norm" (WP 481). DELEUZE AND THE QUESTION OF DESIRE: TOWARD AN IMMANENT THEORY OF ETHICS Daniel W. Smith. **In on a Possible Physical Metatheory of Consciousness Miroljub Dugic et al** show that the modern quantum mechanics, and particularly the theory of decoherence, allows formulating a sort of a physical metatheory of consciousness. Particularly, the analysis of the necessary conditions for the occurrence of decoherence, along with the hypothesis that consciousness bears (more-or-less) well definable physical origin, leads to a wider physical picture naturally involving consciousness. This can be considered as a sort of a psycho-physical parallelism, but on very wide scales bearing some cosmological relevance. Open Systems & Information Dynamics Volume 9, Number 2, 153 (2002) arXiv: quant-ph/0212128. On a broader perspective and background, the theoretical computation of the universe is justified.IN physics and cosmology, digital physics is a collection of theoretical perspectives based on the premise that the universe is, at heart, describable by information, and is therefore computable. Therefore, according to this theory, the universe can be conceived of as either the output of a deterministic or probabilistic computer program, a vast, digital computation device, or mathematically isomorphic to such a device. Digital physics is grounded in one or more of the following hypotheses; listed in order of decreasing strength. The universe or reality: is essentially informational (although not every informational ontology needs to be digital); is essentially computable (the pancomputationalist position); can be described digitally; is in essence digital; is itself a computer (pancomputationalism); is the output of a simulated reality exercise. Following Jaynes and Weizsäcker, the physicist John Archibald Wheeler wrote the following: [...] it is not unreasonable to imagine that

information sits at the core of physics, just as it sits at the core of a computer. (John Archibald Wheeler 1998: 340) It from bit. Otherwise put, every 'it'—every particle, every field of force, even the space-time continuum itself—derives its function, its meaning, its very existence entirely—even if in some contexts indirectly—from the apparatus-elicited answers to yes-or-no questions, binary choices, bits. 'It from bit' symbolizes the idea that every item of the physical world has at bottom—a very deep bottom, in most instances—an immaterial source and explanation; that which we call reality arises in the last analysis from the posing of yes–no questions and the registering of equipment-evoked responses; in short, that all things physical are information-theoretic in origin and that this is a participatory universe. (John Archibald Wheeler 1990: 5) David Chalmers of the Australian National University summarised Wheeler's views as follows: Wheeler (1990) has suggested that information is fundamental to the physics of the universe. According to this 'it from bit' doctrine, the laws of physics can be cast in terms of information, postulating different states that give rise to different effects without actually saying what those states are. It is only their position in an information space that counts. If so, then information is a natural candidate to also play a role in a fundamental theory of consciousness. We are led to a conception of the world on which information is truly fundamental, and on which it has two basic aspects, corresponding to the physical and the phenomenal features of the world. (Wikipedia). **Participatory consciousness forms the bastion, pillar, post, stylobate and fulcrum of the redressal measures and corrective steps taken for circumvention of many problems, identified and unidentified, spread widest commonalty. Anne Baring's sententious and pithy article will be good to have a timely outlook.** I would like to open my contribution to this meeting with the words of Ken Wilber that I quoted in the first Millennium Symposium: 'The secret impulse of life is towards greater consciousness. Maybe the evolutionary sequence is from matter to body, to mind, to soul, to spirit - each transcending and including, each with a greater depth and greater

consciousness and wider embrace' Ken Wilber, and with these from Rupert Sheldrake's contribution to the same meeting: "It is possible to make quite a good case that the sun could be conscious and if the sun, then why not the stars? Why not the galaxies? Why should there not be galactic minds? And if they communicate with each other then the distances are far too great for light to be the medium of transmission. We are into inter-galactic telepathy." After many decades of study, I know of only one diagram in all religious traditions that can adequately explain these statements; that can give us both a blue-print of the many levels of consciousness and a route-map to help us reach beyond our present level of understanding. I would like to introduce a diagram of the four worlds of Kabbalah to illustrate the theme that spirit has brought these four worlds into being and that we, living in the fourth and manifest world know nothing of the existence of the other three invisible ones. It is the clearest diagram I know of that can explain to us the multi-levelled structure of consciousness and our place in the great chain of being. We are indeed, as Wilbur suggests, evolving from matter to body, to mind, to soul, to spirit but where have we come from? It is possible that, as participants in the emanation of spirit into manifest life, we have also come from spirit to soul, to mind, to body, to matter. Everything, from the immense galactic energies to the minute forms of matter participates in one life, is generated from one creative spirit. There is nothing outside this life, this spirit. This diagram suggests that truth is not an article of faith but the experience of participating in the undiscovered reality of what has brought us into being. We cannot know truth until we have entered into communion with that reality. Continuing in similar vein **Michel Bitbol dilates upon the dialectic deliberation and polemical argumentation between** hermeneutists and eliminativists. When he formulated the program of Neurophenomenology, Francisco Varela suggested a balanced methodological dissolution of the “hard problem” of consciousness. He shows that his dissolution is a paradigm which imposes itself onto seemingly opposite views, including materialist approaches. I also point out that

Varela's revolutionary epistemological ideas are gaining wider acceptance as a side effect of a recent controversy between hermeneutists and eliminativists. Finally, he emphasizes a structural parallel between the science of consciousness and the distinctive features of quantum mechanics. This parallel, together with the former convergences, point towards the common origin of the main puzzles of both quantum mechanics and the philosophy of mind: neglect of the constitutive blindspot of objective knowledge. Phenomenology and the Cognitive Sciences Phenomenology and the Cognitive Sciences **Phenomenology and the Cognitive Sciences 2002, Volume 1, Issue 2, pp 181-224 Science as if situation mattered Michel Bitbol**

## ABSOLUTE STABILITY

Hence if $X_8 / M_8^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_8$ ) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8 - 8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2 - 2y_8\frac{(E_8 + n_8^2 y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right) - 4y_8(E_8 - n_8)$$

And has the solution

$$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8 - 8)}[y_8(E_8 + n_8^2 y_8) + [y_8^2 (E_8 + n_8^2 y_8)^2 + 4 y_8 (E_8 - 8)(E_8 - n_8)]^{8/2}]$$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2E_8[y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSOILIDATION**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + (M_8 + U_8)\lambda_8 - (M_8 + V_8)\lambda_8 e^{2T_8\lambda_8} + (M_8 U_8 + Y_8)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8\, e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + (M_8 + V_8) Z_8 \sin(Z_8 T_8) - (U_8 M_8 + Y_8) + (M_8(U_8 + V_8) - X_8)$$

$$\mathrm{X}\cos(Z_8 T_8) - 2M_8 V_8 \cos^2(Z_8 T_8) + M_8 V_8 = 0$$

$$(M_8 + U_8) Z_8 - (M_8 + V_8) Z_8 \cos(Z_8 T_8) + (M_8(U_8 + V_8) - X)$$

$$\mathrm{X}\sin((Z_8 T_8) - M_8 V_8 \sin(Z_8 T_8) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 =$

$1 - \cos Z_8\, T_8$and $A_8 = U_8 - V_8$we obtain

$$4M_8^2 V_8^2 W_8^4 + 2M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8(M_8 + V_8) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2[M_8 - V_8)^2(2A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8(M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8(A_8^2 + 4(X_8 + Y_8))]$$

$$+W_8\left[A_8^2+2X_8\right)\left(2A_8M_8-X_8\right)+2A_8Y_8(M_8+V_8)-M_8A_8^3]+A_8^2\left(X_8M_8-Y_8\right)=0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8>0$$

$$and$$

$$C_8>(X_8-Y_8)/2M_8$$

$recalling\ that\ A_8$= $U_8$and defining $C_8=U_8+V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8(0,U_8V_8=A_8^2(X_8+Y_8)$$

$$Q_8\left(2,U_8V_8\right)=\left(C_8+2M_8^2\right)^2(2C_8M_8+Y_8-X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0$W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$the existence of a D-position at this point is determined by letting cos $(Z_8T_8)=-1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all $T$, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**ABSOLUTE STABILITY IN PREDATOR- PREY MODELS**

**APPLIED TO ACCETUATION ATTRITION MODELS**

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability

criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The Percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending

upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t-T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\big(R(t)\big)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The

regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the

TIME DELAY

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a dominant eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

DELAY DUE TO CONGREGATIONAL AND CONFEDERATIONAL DEBIT CREDIT TYEPES

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t-T_8)F_8(R_8(t-T_8) - N_8(t)M_8((R_8(t) \quad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8} + \frac{N_8 dG_8}{dR_8} - \frac{dB_8}{dR_8} > \frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8} + \frac{dM_8}{dR_8}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_8 = \frac{N_8 G_8}{2M_8}\frac{d(8nF_8)}{dR_8} (\geq 0) \quad (8a)$$

$$I_8 = -\frac{N_8 G_8}{2M_8}\frac{d(8nM_8}{dR_8} (\geq 0) \quad (8b)$$

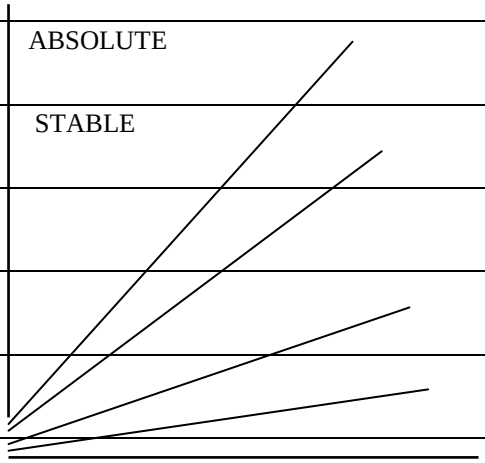
 

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8 > L_8 - I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_8 > L_8 - I_8, \quad \text{when} \quad P_8^2/4 > I \tag{8a}$$

$$P_8 > L_8 - I_8 + \frac{(I_8 - P_8^2}{(L_8 - I)} \quad \text{when} \quad P_8^2/4 < I. \tag{8b}$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit

(like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

See the diagram Nunney gives for predator and prey analysis. We can simulate the same for the financial system in the foregoing too.

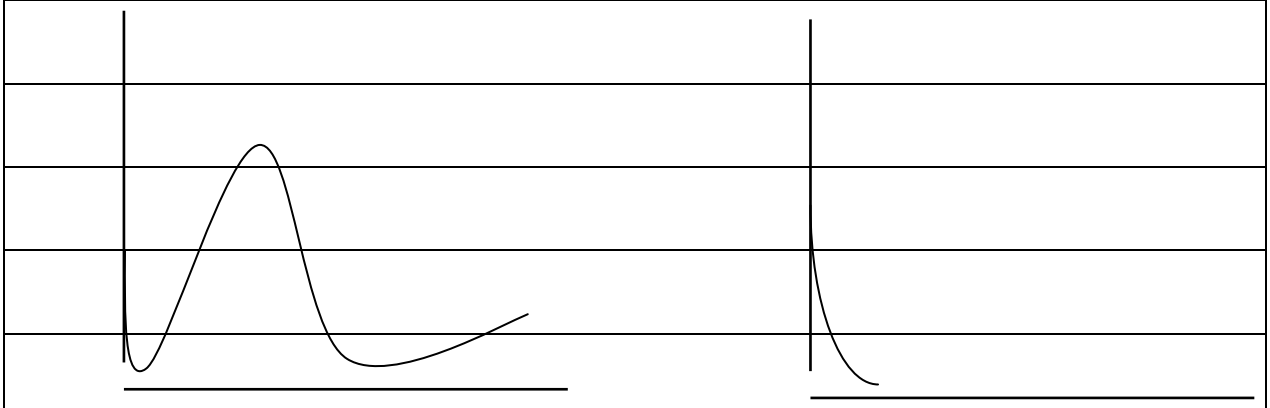
Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the

debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

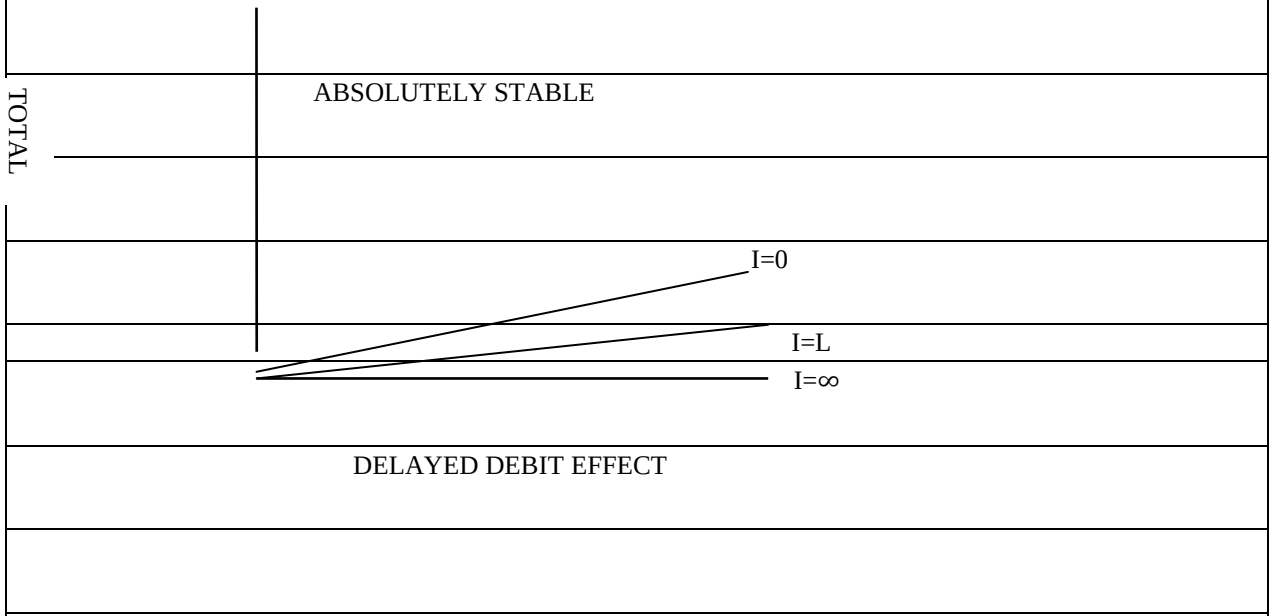
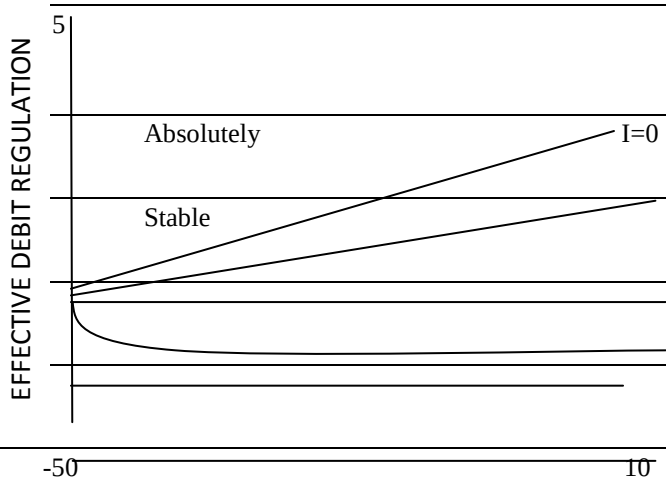



FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison?

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are

defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^\infty \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$ then given the constraints of $f(T_8)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8 \lambda_8 e^{-T_8 \lambda_8} + [(U_8 - V_8) M_8 + Y_8] - [((U_8 -$$

$$V_8 M_8 -] e^{-T_8 \lambda_8} = 0$$

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$\boldsymbol{Z}_8^2 + M_8 Z_8 \sin(Z_8 T_8) - (A_8 M_8 + Y_8) + (A_8 M_8 + X_8)\cos(Z_8 T_8) = 0$$

$$(A_8 M_8) Z_8 - M_8 Z_8 \cos(Z_8 T_8) + (A_8 M_8 - X_8)\sin(Z_8 T_8) = 0$$

Where $A_8 = U_8 - V_8$Using (A2.3), substitute for $\boldsymbol{Z}_8^2$ and $Z_8$in (A2.2):

$W_8^2 (X_8^2 - 3 M_8 A_8 X_8 - M_8^2 (X_8 - Y_8) + 2 M_8^3 A_8 + 2 A_8^2 M_8^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_8 = 1 - \cos(Z_8 T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8 / 2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_8 \le 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2 (X_8 - Y_8)$$

$$Q_8(2, A_8) = (A_8 + 2M_8)^2 (2A_8 M_8 - (X_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2 \left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8 \left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8 \left(2X_8 + Y_8\right)\right)$$

$$+2A_8 \left(X_8 + Y_8\right)$$

$$=2M_8^2 W_8^2 (M_8 + A_8) + M_8 A_8^2 W_8 + (2A_8 M_8 - X_8 + Y_8)(A_8 W_8 + M_8 W_8^2$$

$$+ (2 - W_8)((2M_8 X_8 + M_8 Y_8) W_8 + A_8 (X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $W_8$touches the $W$ –axis only at its minimum value in which case (A2.9) shows that further increases in *A* will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range $0 \le W_8 \le 2$.Defining

$Q_8(W_8, A_8) = a_8 W_8^2 + c_8$ *where* $a_8$ $b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \square b_8 \square 0$$

And

$$b_8^2 \square 4a_8c_8$$

$$Case(i): Y_8 = 0$$

$Let\ Y_8 = 0$ And try the substitution A= $X_8 = /2M_8 in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8 /$ $2M_8$) is positive for *W<2. Since* $\mathrm{Q}_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\ X_8$ ,

$Let\ Y_8 =\ y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 +\ n_8 y_8)$ we have

$$W_8^2 y_8\left[X_8 n_8 \frac{(1 + n_8 y_8)}{(2 + n_8 y_8)} + M_8^2\left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8 X_8 \left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2}+2y_8\left(n_8-1\right)+\frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)}=0.\ (A2.12)\right.$$

Since $n_8$ and $A_8$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8=(2+n_8 y_8),$ *the condition*(A2.88?)*can be written as*

$[\frac{X_8(E_8-8)}{E_8^2}]^2[\frac{X_8^2}{M_8^2}-4y_8 E_8^2\frac{X_8}{M_8^2}+4y_8(n_8-2)E_8^3]$ ?0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8/W_8^2)$ can be expressed as

$2E_8[y_8E_8+(y_8 E_8(8+y_8)(E_8-n_8))^{8/2}]>\frac{X_8}{M_8^2}>2E_8[y_8E_8-y_8E_8(8+y_8)(E_8-n_8))^{8/2}]$

**ABSOLUTE STABILITY**

Hence if $X_8/M_8^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0?$W_8$ ?2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8-8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2-2y_8\frac{(E_8+n_8^2y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right)-4y_8(E_8-n_8)\ ?0$$

And has the solution

$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8-8)}[y_8(E_8 + n_8^2 y_8) + [y_8^2 (E_8 + n_8^2 y_8)^2 + 4 y_8 (E_8 - 8)(E_8 - n_8)]^{8/2}]$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2 E_8 [y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + (M_8 + U_8)\lambda_8 - (M_8 + V_8)\lambda_8 e^{2T_8\lambda_8} + (M_8 U_8 + Y_8)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8\, e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + (M_8 + V_8) Z_8 \sin(Z_8 T_8) - (U_8 M_8 + Y_8) + (M_8(U_8 + V_8) - X_8)$$

$$\mathrm{X}\cos(Z_8 T_8) - 2M_8 V_8 \cos^2(Z_8 T_8) + M_8 V_8 = 0$$

$$(M_8 + U_8) Z_8 - (M_8 + V_8) Z_8 \cos(Z_8 T_8) + (M_8(U_8 + V_8) - X)$$

$$\mathrm{X}\sin((Z_8 T_8) - M_8 V_8 \sin(Z_8 T_8) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 = 1 - \cos Z_8 T_8$and $A_8 = U_8 - V_8$we obtain

$$4M_8^2 V_8^2 W_8^4 + 2M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8(M_8 + V_8) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2[M_8 - V_8)^2(2A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8(M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8(A_8^2 + 4(X_8 + Y_8))]$$

$$+ W_8[A_8^2 + 2X_8)(2A_8 M_8 - X_8) + 2A_8 Y_8(M_8 + V_8) - M_8 A_8^3] + A_8^2(X_8 M_8 - Y_8) = 0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8 > 0$$

$$and$$

$$C_8 > (X_8 - Y_8)/2M_8$$

$recalling\ that\ A_8 = U_8$ and defining $C_8 = U_8 + V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8(0, U_8 V_8 = A_8^2 (X_8 + Y_8)$$

$$Q_8\left(2, U_8 V_8\right) = \left(C_8 + 2M_8^2\right)^2 (2C_8 M_8 + Y_8 - X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$the existence of a D-position at this point is determined by letting cos $(Z_8 T_8) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all *T*, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**ABSOLUTE STABILITY IN PREDATOR- PREY MODELS**

**APPLIED TO ACCENTUATION ATTRITION MODELS**

We apply model to each and every module and a concatenated model was to be presented which we procrastinate due to spatial restrictions.

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - D(R) - NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

The question is political-discursive. It's what Stengers and Latour are getting at with their concepts of diplomacy and cosmopolitics. They grant, first, that all entities exist and, second, that to say that someone's cherished idol (or whatever disputed entity they hold dear) is non-existent is a **'declaration of war'** – 'this means war,' as Stengers often says. They thus shunt onto-political discourse off of the terrain of knowledge/belief in the sense of existence/non-existence. Their basic claim seems to be that 'respect for otherness,' i.e. political pluralism, can only come from granting the entities that others hold dear an ontology, even if you don't 'believe' in them. You

are thus permitted to say 'I do not follow that god, he has no hold over me' but you are not permitted to say 'your god is an inane, infantile, non-existent fantasy, grow up.' And it's not just a question of politeness (although there's that too). The point is to grant others' idols and deities an existence – one needn't agree over what that existence entails, over what capacities that entity has or what obligations it impresses upon you as someone in its partial presence but to deny it existence entirely is to 'declare war' – to deny the possibility of civil discourse, of pluralistic co-existence. I believe that it was Richard Rorty who once quipped something like claims to reconcile realism and idealism always seem to end with a triumph of idealism. We don't, in fact, get realism through such approaches, but rather just get a pervasive anti-realism. I think this is also the problem with the "non-controversial pluralism" advocated by Stengers and Latour that Phillip defends here. Such pluralism is not realism but is, in fact, a thoroughgoing social constructivism. I think this is the central problem with Latour's argument in Irreductions (these days I regret having ever defended it). In rejecting both Enlightenment critique and what he calls "reduction" he wants to say something like "The Pentacostal really is filled with (e&eb) the Holy Spirit", that for the 19th working scientist heat really is a fluid and phlogiston really is what allows things to burn, and that for the Greek lightning really is an expression of Zeus's anger. Latour tells us that we aren't to reduce or explain away the entities posited by another group's "ontology" but are to develop explanations from within that ontology. It might not sound particularly sexy– and it certainly doesn't tell us what is worth thinking –but I can't help but believe that philosophy is the critical and reflective investigation of basic concepts that guide our investigation of the world about us, how we ought to live our lives, and what form of governance might be best. Compare two figures. A scientist might ask, what causes depression? We can very well imagine a philosopher turning around and asking the scientist, what is causality? The scientist presupposes a concept of causality in her investigations. She uses this concept in her inquiry. Now she might have a sophisticated concept of causality or she might never have thought much about causality at all, using it in the sort of colloquial and unreflective way that Plato decried when, for example, people

like Euthyphro talked about piety. A whole cascade of questions arise when we raise a question like "what is causality?" We can ask whether or not causality exists at all. We can ask how we distinguish between correlation and two events that merely accompany one another from genuine causation. This, for example, was Hume's question. But perhaps most importantly we can ask whether there is only one form of causality or many forms of causality. Is there only one-to-one causation; one cause and one effect? Is there many-to-one causation; or many events conspiring to produce an effect? Is there one-to-many causation; or one event producing a variety of different effects? We can even ask whether causality necessarily moves from past to present or whether there aren't forms of causality that move from future to past! **January 23, 2014 Pluralism and Realism:** multiple-worlds Over at Struggles Forever, Jeremy Trombley has an interesting post up on "the ontological turn" in anthropology or ethnography. I've been meaning to have a discussion with him about this as I think it's an issue many of us are struggling with. For example, the core project of The Democracy of Objects– a project which I think many have missed –is to somehow reconcile some version of social constructivism with a realist ontology capable of making room for ecology (which requires realist and materialist positions as there's a fact of the matter where global warming is concerned) as well as the role played by objective agencies in social assemblages such as technologies, infrastructure, features of geography, local climates, the growth cycles of plants and animals, waste, etc. Maybe we can try to organize some cross-blog event to discuss these issues. I certainly think they're close to the heart of Jeremy, Michael of Archive Fire, Arran James, and a host of others. As an aside, I'm beginning to realize how the different sites of the political I've been outlining– semiopolitics, thermopolitics, oikopolitics (political economy), geopolitics, eropolitics (the politics of sex and desire), biopolitics, and chronopolitics (and I'm sure there are other political sites!) –are drawing me away from traditional Marxism. Assuming that classical Marxism holds that economics or the conditions and relations of production are determinative of all other sites of the political, the various sites of the political that I've been outlining would lead to the conclusion that there is not one determinative

base of the political. This would not require committing Marx to flames, but rather of recognizing the phenomenon of overdetermination, or of a variety of different entangled sites of the political. But I digress. First, I find myself wondering what the ontological turn means in ethnography. Is it 1) the investigation of the different ontologies held/proposed by different cultures? E.g., the Aztecs believed that reality was structured in this way, while the Greeks in that way, and the ancient Chinese this way, etc? Or 2) is it an investigation of how real entities– independent of cultural beliefs –influence cultural formations? Or is it a combination of both? A position that I would favor. January 20, 2014 Model is applied to structured reality

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

**The evaporation spectrum of black holes from a local quantum gravity perspective Aurélien Barrau (Submitted on 21 May 2016**)

(9) Authors revisit the hypothesis of (e) a possible line structure in the Hawking evaporation spectrum of black holes, due to (e) non-perturbative quantum gravity effects, even arbitrarily far away from (e) the Planck mass.

(10) Authors show that this naive prediction might in fact hold in (eb)the specific context of loop quantum gravity, with (e&eb) a small departure from (e) the ideal case for (e) some low-spin transitions.

They also show that the effect is (=) neither washed out by the dynamics of (e) the process, nor by (e) existence of a mass spectrum up to (e) a given width, nor by the secondary component induced by (e) the decay of neutral pions emitted during (e&eb) the time-integrated evaporation. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Cite as: arXiv:1605.06613 [gr-qc] (or arXiv:1605.06613v1 [gr-qc] for this version

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t - T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\left(R(t)\right)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

It's very easy to treat this as an abstract, academic debate: "Are you a realist or are you an anti-realist?", as if it were just a matter of what happens to be true. But it seems to me that this debate has, in reality, always been about politics. As I outlined in my last post, we have perpetually seen how appeals to the real and natural have been used in the name of oppressive power, inscribing both the exploitation of nature and the oppression of various people in the very fabric of being itself. Theistic theology and realist ontology have perpetually been used in the name of

what Deleuze called "State Philosophies" or philosophies that ontologize contingent orders of power and privilege (e.g., "the great chain of being" used to justify patriarchy, monarchy, serfdom, poverty, etc, and appeals to nature used to justify poverty and racial inequality (The Bell Curve), patriarchy (evolutionary psychology), heteronormativity, capitalism, etc). Because arrangements of power and inequality are always contingent in the sense that there's no marked difference in the capacities of peoples, power always looks for a transcendent supplement that would provide justification through ontological necessity. Antirealism– from the Greek atomists to present –became the radical and emancipatory gesture because it revealed the lie behind all of these forms of social organization or their inherit contingency or arbitrariness. Realism, by contrast, has all too often functioned as an apologetics for arbitrary power and social organizations. Here it's worth recalling what Foucault said about science: "…Even before we know to what extent something like Marxism or psychoanalysis is analogous to a scientific practice in its day-to-day operations, in its rules of construction, in the concepts it uses, we should be asking the question, asking ourselves about the aspiration to power that is inherent in the claim to being a science. The question or questions that have to be asked are: "What types of knowledge are you trying to disqualify when you say that you are a science? What speaking subject, what discursive subject, what subject of experience and knowledge are you trying to minorize when you say 'I speak this discourse, I am speaking a scientific discourse, and I am a scientist.' What theoretico-political vanguard are you trying to put on the throne in order to detach it from the massive, circulating, and discontinuous forms that knowledge can take?" (Society Must Be Defended, 10). All of these questions hold equally for claims to something being real. What is one trying to minorize when claiming something is real? What becomes privileged? What is excluded? It is these questions that have been at the heart of the realism debates, for as Spencer-Brown taught us, every distinction has a marked and unmarked space, draws attention to something to be included and pushes something into the unconscious or the domain of the invisible, hidden, or veiled. This is above all the case with evocations of the real. However, as I've tried to show antirealism

leads to its own problems. First, so long as we exclude real beings from our ontological inventory, we are unable to fully understand how power functions ("such and such a set of cultural formations have been thoroughly debunked, yet people still live as if they believed them"). Not only do we not fully understand the sources of the problems due to too much focus on the discursive and semiotic, but we deny ourselves valuable sites of political intervention at the level of infrastructure. Second, we are prevented from addressing things such as cultural racism, such as that found in Heidegger's privileging of the West and the Greeks and Germans in particular. We do a good job addressing biological and theological racism, heteronormativity, and sexism by showing how it is a cultural construction, but when faced with racism such as Heidegger's where he argues that there's something "unique" about the Greek event, it's language, and about the German language, or Badiou/Zizek's racism with respect to the "Pauline Event", we really have no response. Here someone like Jared Diamond or Fernand Braud el is needed to explain global-geographical inequalities. Third, the tools of the cultural turn really do not provide us with the means of thinking the ecological as a site of the political. For this reason, I've tried to formulate a third way– which might be called "constructivist realism" or "constructivist naturalism" –that retains the insights of the cultural turn, while also allowing a robust place for the material. I don't claim to be original in this. I think that many such as Manuel DeLanda, Deleuze and Guattari, Stacy Alaimo, Karen Barad, Jane Bennett, etc., are up to something similar. It's a vast project that requires the work of a multitude of voices, especially given the way in which the culturalism pervades humanities. **The Relational and the Non-Relational: Notes towards an Immanent and Pluralist Theory of Meaning Posted by larval subjects under uncategorized k-bigpic I have deleted or reformulated some sentences for which I beg pardon attributable to spatial constraints** What is correlation, what is causation? What is coordination? What is coincidence? These thoughts have bothered humans since long. I have been able to excavate some reliable material from Wikipedia which I would like to share. "All views are partial and contingent – that's the lesson of pluralism, and you say as much yourself when you say that we must be attentive to ways our own

knowledge might contain superstitions. But it's not discourse that determines whether one view works and another doesn't – it's a confrontation with **non-discursive (non-human)** agencies." I wonder if part of the issue here is that we understand pluralism differently. For me this doesn't sound like pluralism at all. Rather, it just sounds like our **epistemological condition**. We [hopefully] come to understand those portions of existence we question or investigate. The reason we come to investigate them is largely contingent. Finally, the accounts of these features of existence we give can be mistaken. All of this is perfectly consistent with a monism. I'm not sure why what he outlines above is the lesson of pluralism. It seems to me that every realist knows this. It seems to me that ontological pluralism is something quite different. Ontological pluralism is the thesis that there are many different worlds inhabited by many different entities. Thus, for example, you would have one world; say that of Lucretius, that's only inhabited by atoms and their combinations. You would have another world; say the world of the Mongolian shaman, that's inhabited by spirits that do all sorts of things. The ontological pluralist is saying that these spirits are, that they exist, that they're real. This is something quite different than merely saying that there are partial and contingent points of view on the world. It's this saying that these entities are rather than that some person or groups of people believe that they are that is the nub of the issue. Now, I think part of the issue here is that there's an ambiguity in the term ontology. Ontology can be one of two things. On the one hand, ontology is a group or persons set of beliefs as to what is. Here it's trivially true that there are pluralities of ontologies and the realist readily recognizes this. This is the whole reason there are debates over ontology. Mongolian shamans have their ontology, Europeans theirs, Christian fundamentalists theirs, materialists theirs, etc. When striving to understand and communicate with others it's vital to understand these ontologies because, as rhetoricians like Burke point out, our **beliefs about what is** are among the things that motivate our action. Despite having never seen bacteria I was my hands because I believe there are bacteria and viruses on door handles and whatnot and don't want to get sick. My belief about a particular thing existing is what motivates my action. And who knows, perhaps this belief is as

superstitious as the belief that the crops failed because God was displeased with my community. On the other hand, ontology is a theory **about what is.** It is making a claim that something exists. This is where the rubber hits the road. The ontological pluralist seems committed to the thesis that every group set of beliefs about what exists is sufficient for granting the existence of those entities. This is what I find objectionable in Latour. Obviously I'm not bothered by Latour's suggestion that we should take into account the role that nonhumans like speed bumps, rivers, microbes, etc., play in the form that social assemblages take. But this is not all that Latour claims. This is precisely where the philosopher might balk. Unlike the ethnographer, the philosopher is not interested in what people believe exists, but rather philosophers– at least of the realist variant –are trying to figure out what is. In other words, the realist philosopher begins with the premise that not all of these beliefs about what is are true. So for the philosopher, recognizing that Mongolian shaman's believe in the existence of shamans would only be the first step. The next step would consist in determining whether there's good reason for thinking such entities really do exist, i.e., whether there's good reason for believing these entities have mind and culture independent reality. Lest readers think that I'm just picking on the supernatural here, we can ask similar questions about strings, subatomic particles, galaxies, etc., etc., etc. I guess the question really comes down to what exactly we mean by ontological pluralism. When we talk about ontological pluralism are we defending the thesis that people have different theories of being? If so, then ontological pluralism is trivially true. If this is what is meant, then I certainly share Jeremy's view that it's valuable to understand the different world's people believe in. Certainly when I was practicing as an analyst I didn't get in ontological debates with my patients and it was necessary to understand their theory of being or their ontology to properly attend to them. Or, when we talk about ontological pluralism are we defending the thesis that all these ontological theories are true and refer to really existing entities? That's quite a different claim and is not one I would defend or endorse. Now someone might object that "in both The Democracy of Objects and Onto-Cartography you defend the thesis that there are multiple worlds." This is true. Because I hold that not everything

is related I'm led to the conclusion that there are diverse worlds. However, I also hold that however many worlds there might be, these worlds are nonetheless composed solely of material entities. Within the framework I propose I wouldn't suggest that there's one world where there are spirits and another world where there are souls and yet another where there are only material entities. My view is that there aren't spirits or souls in any of these worlds. One of the things I keep hearing in these discussions is that somehow the realist adopts a view from nowhere. I honestly don't understand this criticism. Investigation always occurs somewhere and requires all sorts of mediations involving technologies, experiments, etc. It's that labor of gathering evidence, conducting experiments, using technologies to observe the world, etc., that gradually gives us a body of data that allows us to say there's good reason to believe that such and such a thing exists and has these powers. Another charge seems to be that the realist refuses to recognize that their claims about the world are fallible. I find this charge particularly strange because it's precisely because the realist recognizes the difference between our theories of the world or what we say about the world and the world itself that fallibility is built into the core of his position. Realism doesn't mean one holds they have special access to the world, that they know all truths, or that they have the truth in hand, only that there are truths to be known and that we can be mistaken about things. Circling **Squares Philip** has a post responding to my quandaries about how to mesh realism and pluralism. He writes: Ontologically and metaphysically the idea of realist pluralism is no longer an issue. There are (appropriately) numerous variants but the basic idea that reality is itself pluralistic is well established. Captioned expatiation is about Everett worlds and string theory

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from

being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

The von Neumann formulation leads us to conclude that the lower realms of quantum reality form a virtually limitless array of potentials that are inherently incapable of realization without the observation of a knowing entity, which we identify with consciousness (Kafatos and Nadeau, 1990). We thus have a correspondence between quantum theory, the hierarchical structure implied by the Holographic Principle, and the higher orders of experience,

leading to the full expression of the Conscious Universe. David Bohm's theory of the implicate and explicate orders involves a holographic Principle that is fully consistent with the Holographic Principle discussed here (Bohm 1980; 1986; Bohm and Hiley, 1993). According to Bohm, there is an implicate order that represents the universal, holographic subtext of reality, and which unfolds in every moment to produce the explicate order that we all observe. Thomas Germinario (2004) has equated the implicate order with (eb) the unconscious process, and the explicate order with (eb) conscious process. He emphasizes the nature of dreams within (eb) the implicate order, and the importance of the dream work for maintaining a healthy mind through (e&eb) integration of the implicate order and our daily lives. Allan Combs and Mark Holland (1990) connected the implicate order or holomovement (Bohm, 1980) with (e&eb) Carl Jung is theory of synchronicity (Jung and Pauli, 1955), with (e&eb) the implicate order providing (eb) a holographic medium through which apparently disconnected individuals become connected. The principle of synchronicity, the instantaneous connection of people and events beyond (e) the senses, has been equated with the quantum-physical principle of non-locality (Combs and Holland, 1990; Germine, 1991), and has been proposed to be the fundamental mechanism of conscious process (Germine, 1991).Applying the Holographic Principle theory of mind, subjects may be connected synchronously in the manner of two individuals conversing through cellular phones, where the radio signals are transmitted through (e&eb) a distant satellite. The difference here is that the radio signals travel at the speed of light, and thus the there is some miniscule time lapse between the sending and receiving ends. The Holographic Principle implies (eb) a much more distant and instantaneous communication pathway, with (e&eb) the analogue of the satellite being the holographic boundary of the Universe, as well as a much richer communication, potentially involving a transmission of thoughts and feelings. Our ego-consciousness seems to mask (e) the universal relatedness implied by (e) the Holographic Principle, and it is perhaps only through transcendence of (e) the ego-consciousness that the higher orders of experience can become conscious. In the phenomenon of synchronicity, there seems to be a meaningful connection between (e&eb) individuals that

breaks through the barrier of ego-consciousness. Such a connection is reported by many individuals in the course of dreams, when (e) the ego-consciousness has been suspended, at times informing (e&eb) the dreamer of something that has happened in the life of a meaningfully-connected individual. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Impassioned-intense Connectionist neuroscience, in an effort to attain relevance to clinical practice, has recently given us a theory of the relationship of the unconscious to (e&eb) consciousness which is (=) dualistic (Viamontes and Beitman, 2007), an ad hoc theory of conclusions based on (e) assumptions, the most erroneous of which are that representations of (e) information are information, and that such representations elaborate themselves through (e&eb) a process of representations of representations in neuronal circuits (Bechtel and Abrahamsen,1991). As theory, the connectionist models are lacking (e) heuristic value and **are (=) inimical to the progress** of (e) basic and clinical neuroscience. This is not to denigrate or minimize the importance of neuronal connections, but these connections are rather like the wires in (eb) a radio that plays a symphony. The radio, in and of itself, is (=) incapable of playing (e&eb) the symphony. The radio doesn't do not write the symphony, nor does it broadcast the Passionate-raging symphony over the airways. Yet, if we loosen one wire, the symphony is (=) not manifested through the radio. The brain is (=) not subject to (e) the same vulnerability, in that it has (e) parallel and redundant networks of processing, and, it is in this sense that the connectivity of the brain becomes (eb) important. Connectionism, to the extent that it accepts (e) consciousness and the (e&eb) unconscious at all, attributes them to (e) separate groups of "circuits" in the brain, with (e&eb) consciousness-processing circuits having (e) a limited but detailed ability to analyze (e&eb) information, and unconscious-processing processing more information in less detail (Viamontes and Beitman, 2007). There is an element of truth in this view in that information that has importance and meaning resonates within the brain at higher levels of experience and at higher levels of holographic recursion. This resonance involves what we know, what we feel, and what we are capable and willing to consciously realize. The Universe, as William Blake noted,

can be seen in a grain of sand. Indeed, the grain of sand requires (e) the Universe to exist. The human organism requires (e) the Self to exist. The Self is not a simple amalgam of (e) sensory experiences, as suggested in the connectionist model (Blinder, 2007). This amalgam view of the Self can be damaging when applied to (e&eb) psychotherapy, as the connectionists suggest (Blinder, 2007). The Self is ours forever and ever, from eternity to eternity. As Milton suggested, the mind is (=) its own place. Our deepest purpose is for our minds to resonate with (e&eb) the supra-conscious levels of experience, much as it is (=) purpose of the radio to play the symphony. If the symphony is not to our liking, (eb) we can turn the radio off. So, perhaps, the transpersonal consciousness can turn (e&eb) our own individual and collective radios off, or perhaps turn the volume down, if (e) we are not in harmony with the **Universal Mind**. Physics has given us many enigmas. Quantum theory itself is an enigma. But some of the enigmas seem to involve (e&eb) the limitation of (e) our own abilities of discernment. The thought experiment of (e) "Schrödinger's Cat," as it is often interpreted (e.g. Gribbin, 1984), assumes (eb) that the **cat is not conscious, and is (=) incapable of observation in the context of (e) collapse of (e) the wave function.** This is due, in large part, to the false dichotomy of (e) information and (e&eb) experience. Our observation of other animals indicates, to those that are attuned to (e) the animal mind, who can "feel" the animal mind, that, at (eb) the very least, all mammals are (=) conscious. It would seem that the great neuropsychiatrist, Stanley Cobb (1948), may have been correct in his attribution of (e) levels of consciousness to a wide variety of organisms: "...lower animals with (e&eb) no cerebrum appear to be conscious...even plants such as the sunflower that turns towards strong light may have (e) a vague awareness or warmth and comfort. There are many degrees of consciousness and it is my contention that it is integrated at (e&eb) many levels like other important functions of (e) the central nervous system."With respect to the theory of mind, recent mainstream thinking remains (eb) classical and mechanistic. Such views of mind are advocated by our most in prominent neuroscientists (e.g. Changeux, 1985; Gazzaniga, 1985) as neuroscience research attempts to explain all mental function on the basis of (e) **pure mechanism and the localized function of**

**specific areas of the brain.** Francis Crick, the distinguished Nobel Laureate, had a very influential second career as a neuroscientist. He concluded that we are basically soulless creatures, that all mental processes could be reduced to (e&eb) identifiable, neural correlates in (eb) the brain, and that, in fact, we are (=) nothing but a collection of neurons (Crick, 1995). Crick even went so far as to say that what we call the soul is (=) a group of neurons located in the prefrontal lobes of the brain. Another Nobel Laureate, Gerald Edelman (1987), has written a critically acclaimed book on "neural Darwinism," which purports that the brain is "circuitry" is constructed through a developmental process of "survival of the fittest" neurons and neuronal connections. At the same time the total purposelessness and lack of direction in evolution is expounded, again to critical acclaim, by Richard Dawkins (1986). The idea that there essentially no self or soul imbued with agency and self determination was expounded by Gilbert Ryle (1949/2002) with his parody of the "ghost in the machine," and this work continues to be influential among many scholars of the mind. Philosophers such as Patricia Churchland (1986), who coined the term "neurophilosophy," take an "eliminative materialist" view of (e) belief, free will, and consciousness. Other philosophers, such as Daniel Dennett (1991) would make (eb) experience "epiphenomenal," a by-product of (e) brain processes with (e&eb) no effect whatsoever on (eb) the function of mind. However, unlike neuroscience, which is still in its infancy, the philosophy of mind has had a long history, including such notable figures as Plato, Aristotle, Descartes, Kant, Locke, Hume, Kierkegaard, and James, to name just a few. Perhaps, as noted by Ervin Laszlo (1974): "in today's world, most of the traditional functions of (e) cognitive synthesis have atrophied (e) and are ignored and neglected." Hopefully, what we have presented here will be a step in the right direction for what Laszlo calls a "conceptual synthesis" to help fill the need for (e) meaningful engagement is such a world. The fundamental elements of such a synthesis have been described by (e) Laszlo as follows: Conceptual synthesis performs (e&eb) at least five basic functions in the guidance of affairs. They are the mystical, the cosmological, the sociological, the pedagogical or psychological, and the editorial functions. The mystical function inspires (eb)

in man a sense of mystery and profound meaning related to (e&eb) the universe and of himself in it. The cosmological function forms (eb) images of the universe in accord with (e&eb) local knowledge and experience, enabling (eb) men to describe and identify (e&eb) the structure of (e) the universe and the forces of nature. The sociological function validates supports and enforces (eb) social order, representing (eb) it in accord with (e&eb) the nature of the universe, or as the natural or right form of (e) social organization. The pedagogical or psychological function guides (eb) individuals through stages of life, teaching ways of understanding themselves and others and presenting (eb) desirable responses to life is challenges and trials. Finally, the editorial function of (e) conceptual is to define (eb)some aspects of reality as important and credible and hence to be attended to, and other aspects unworthy of serious attention. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Kant sums up this inversion and its spirit early in the Critique of Pure Reason: Up to now it has been assumed that all our cognition must conform to the objects; but all attempts to find something about them a priori through concepts that would extend our cognition have, on this presupposition, come to nothing. Hence let us once try whether we do not get farther with the problems of metaphysics by assuming that the objects must conform to our cognition, which would agree better with the requested possibility of an a priori cognition of them, which is to establish something about objects before they are given to us. This would be just like the first thoughts of Copernicus, who, when he did not make good progress in the explanation of the celestial motions if he assumed that the entire celestial host revolves around the observer, tried to see if he might not have greater success if he made the observer revolve and left the stars at rest. (B xvi) Onticology, like all variations of object-oriented ontology, is realist in its orientation. In defending a realist ontology onticology holds that the vast majorities of objects, actants, beings, or entities are independent of humans and are what they are regardless of whether any humans regard them or register them. In short, onticology rejects any anthropomorphic, idealist, or anti-realist thesis to the effect that to be is to be the correlate of mind, spirit, the body, the human, and language or otherwise. While it is certainly the case that knowledge is necessarily dependent on the object to which it relates, the converse does not hold true. Objects are not dependent on being known, regarded, perceived, or spoken

about. As such, and to put it in Aristotlean terms, knowledge is an accident of objects, not objects an accident of knowledge. As Althusser so nicely puts it, "[n]o doubt there is a relation between thought-about-the-real and thisreal, but it is a relation of knowledge, a relation of adequacy or inadequacy of knowledge, not a real relation, meaning by this a relation inscribed in that real of which the thought is the (adequate or inadequate) knowledge" (Reading Capital, 96). Althusser goes on to remark that "[t]he distinction between a relation of knowledge and a relation of the real is a fundamental one: if we did not respect it we should fall irreversibly into either speculative or empiricist idealism" (ibid.). Onticology categorically endorses Althusser's verdict. It is a fundamental necessity to distinguish between those relations that belong to the object and those that belong to knowledge. Contemporary philosophy continuously confuses these two very different sorts of relations. Naturally the question arises of how it is possible to surmount our relation to the object so as to determine whether objects themselves possess the properties we encounter in relating to objects. In other words, given that we can only ever relate to the object in relating to the object how is it possible to surmount this relation to get at the being of the object itself? Much more will have to be said about this later– and the answers will be surprising with respect to standard prejudices about realism –however, for the moment it can be said that onticology takes its epistemological inspiration from the transcendental realism of Roy Bhaskar. Among other things, Bhaskar sought to provide a transcendental grounding for the sciences. Insofar as onticology defends the thesis that the field of being is much more vast than the field of objects investigated by the natural sciences, it parts way with the thesis that the domain of being is exhausted by the domain of natural objects. However, the general form of Bhaskar's argument holds for our realist purposes. A transcendental argument seeks to elucidate the conditions under which certain acknowledged practices and forms of cognition are possible. Kant, for example, asked what must be the case for mathematical judgments to be possible. How is it both that we are able to extend our knowledge, as if by magic, through mathematical judgments and, more significantly, that these judgments are able to provide genuine knowledge of the world despite the fact that these

forms of reasoning are not based on experience? Part of Kant's argument consisted in claiming that mind imposes the forms of space and time on the data of experience. In other words, space and time are not attributes of being itself but rather of the mind that regards being. Insofar, Kant argues, as mathematics is ultimately a rumination on the nature of space and time taken in their most abstract form and insofar as the mind imposes space and time on the manifold of sensation, it thus follows that a priori judgments about the nature of spatio-temporal relations are possible that anticipate the structure of actual-space times without directly experiencing these space-times. Why? Because any manifold of sensation must necessarily be structured by these forms imposed by intuition. **Onticology–A Manifesto for Object-Oriented Ontology Part I Posted by larval subjects under Object-Oriented Philosophy My Question Is This: You Say, "Granted There Are Deeper And Deeper And More Primitive Forms Of "Knowledge" Or Abstraction Into Mechanism, But I Wouldn't Say That The Atomic Level Brings A New Epistemic World Into View." But Doesn't Any Kind Of Knowledge, However Primitive, Imply Bringing Some Kind Of Epistemic World Into View Or Proto-View?** Joel: David, Yes Exactly. My Question Is, Just How Far Down The Complexity Gradient Can The Meaning Of 'Knowledge,' Or 'Semiotics, Or Specifically Representation (Maya) Be Stretched Without Breaking? I Can Comfortably Stretch It To The Genetic Level, As The Proto-Epistemic, Where Genetic Codes 'Represent' Phenotypic Structures Or Modes, And The Whole Of Evolution Learns From Its Mistakes And Successes As That Knowledge And History Is Encoded Into And As The Recapitulation Of The Embryogenesis Of The Organism. But An Atom, In My Model, While Certainly Reacting To Its Environment From Its Own Attractor, Essence, And 'Interiority' (Prehension) Isn't Making Representations Of It To Do So. But The Key To The Power And 'Essence' Of Representation (And Simultaneously Its Limitation) Is The Capacity Of Abstraction And Choice. Evolution Does This At A Rudimentary Level (And Anthropomorphizing A Bit) With Its Sending Into The World Its 'Random' Variations, Organismic Ideas, In A Sense. Its 'Mind' Is Literally On The Outside As Our Living World And Biosphere. A 'Choice' Is Made In The Biospheric Interiority,

And Knowledge For The Evolutionary/Embryogenetic Trajectory Is Gained By The "Differential Reproductive Success" Of These Individuals Lives Through Time, As The Increasing Intelligence, Representation, And Capacity For Choice Is Directly Injected Into The Flow Through Sexual Selection. But Where Is The Encoded/Abstracted (Enfolded) Representation Of Options And A Choice Among Them Made At The Atomic Level? I Am Working With The Most Cutting-Edge And Coherent Models For The Atom (Including My Own) That I Can Find Among The Heterodoxy, And I See No Evidence Of Representation And Choice At This Level, Or Rather No Real Way To Further Stretch The Meaning Of The Terms To This Level. Sure There Is Uncertainty From Infinite Difference And Immanent Causation (See Bohm On Infinite Causation), But Without Representation There Really Is No Choice (Although Crudely We Could Call A 'Bifurcation' A 'Choice' By The System As A Whole). So I Am Simply Saying That Representation At The Atomic Level Is Entirely Enfolded And These Are Really Just Ontic Level Phenomena. The Atomic Level, In My View, Is Pre- Not Yet Proto-Epistemic, Which I Place at the Genetic Level Where We Find the First Codes And (e&eb) Primitive Representation. This Just Means That There Is A Real Distinction Between Ontic And Epistemic (And A Gradient Between Them) And We Can't Collapse The Two Into Just The Ontic-Epistemic. And Yet, With The Distinction Intact, And The Gradient To Explore, They Remain Nondual, Univocal, ONE Or Ontic. What I Would Say Exists At All Levels (My Own Self-Aware Myth About The Given) Is What I Call The "Symbiogenesis Of Subject And Object". This Is The Spinozan Essence Of Dynamic Stability And Growth, And Leibnizian Prehension. Also Called The "Nucleation Of Observability" In Spinbitz. It Is From This Nucleus Of The Observer (Fuller) In A "Point-Free Geometry" (Whitehead) Of Pointless "Points Of View" As Conditions Of Boundary That Wilber's "All Is Perspective" Finds Its Ontic Grounding. And This Symbiogenesis Is Also A Rudimentary Or Deep-Level Native Or Primitive Intelligence, Evolution At The Involution Of The Subject-Object Interface. It Is An Exploration Of Creative Learning And Direct Immanent Awareness, Unmediated (Enfolded) By Representational Forms (Unfolded). It Just Is, In Waves (Cosmic Vertebrae)

Before It Abstracts, Represents And Incarnates The Flesh Of What Could Be. Matter Is Enfolded Seed Of The Flower Of Its Abstraction. Maya Is This Depth And Recursion Of Brahma Into And Through The Boundary Conditions Of Abstract Relation. And Brahma Is The Span Of The Emergent Rigor Of The Cosmic Ergodic Spine That Opens The Recursions Into The Field Of Infinite Difference. In Immanent-Transcendent Waves, Literally And Empirically A Limit-Cycle In Ergodicity, It Enfolds And Unfolds From Anatom To Anatomy, From Infinite Intelligence Or Omnirelational Awareness Of Self, Through The Representational Depths Of Self As Self-Consciousness. Omni-Evolution, Trans-Dynamics, And The Anatom Posted On February 21, 2013A Discussion With Tom Huston, Part II The Following Is Part 2 Of An Excerpt From A Discussion Initiated By Tom Huston On His Facebook Page, Edited And Embellished For Clarity. Find Part One, Here. In This Part We Dig Into The Nature Of The "Anatom" And The "Ergodic Spine" Of An Electro-Fractal Cosmos, And Into The Trans-Dynamic Integration Between Being And Becoming, And The X-Interface−The Crossroads Of The Ontic-Epistemic (Brahma And Maya) And The Subject-Object Polarities. Univocal_Dynamics_V1 Tom Huston Is A Founding Member Of Integral Institute And A Former Editor Of Enlighten Next Magazine. His Webpage Is Tomhuston.Com. David Marshall Is A Writer And Editor Living In Chicago. No meaning, value, or normativity to be found in nature, that there is nothing **natural beings ought to be,** but that, rather, these judgments arise from us (See last post of larval subjects) which has led some to raise valuable questions about the coherence of these claims. The problem arises when the following three propositions are taken together: 1. there is nothing outside of nature.2. Beings have no intrinsic meaning, purpose, or value (in and of themselves, there's nothing they ought to be). 3. Value judgments about what beings and being ought to be arise from us and beings like us (bonobo apes, dolphins, institutions, birds of paradise, etc). The problem arises between proposition 1 and 3. How can it both be true that there is nothing outside of being and those normative judgments belong to us and other beings capable of making normative judgments, not nature? The problem arises from restricting these judgments to humans and beings capable of

making these judgments. In making such a claim it seems as if we're saying that there's something outside of nature, something that is beyond nature, thereby violating the first thesis and potentially reintroducing the nature/culture distinction. By contrast, if we say that normative judgments are the special domain of those living beings with the proper degree of sentience to make such judgments, then we seem to reintroduce the nature/culture distinction and fall back into the sorts of problems that I outlined in my last post (and that are so nicely critiqued by thinkers such as Latour). Is there a way out of this? I don't know. Great_Chain_of_Being_2_ (lighter) ASIDE: It seems to me that this is really what the debate between realism and anti-realism, realism and socio-linguistic constructivism surrounding the new materialisms and speculative realists has really been all about.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \qquad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9}+\frac{N_9 dG_9}{dR_9}-\frac{dB_9}{dR_9}>\frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9}+\frac{dM_9}{dR_9}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usance bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_9=\frac{N_9 G_9}{2M_9}\frac{d(9nF_9)}{dR_9}(\geq 0) \qquad (9a)$$

$$I_9=-\frac{N_9 G_9}{2M_9}\frac{d(9nM_9}{dR_9}(\geq 0) \qquad (9b)$$

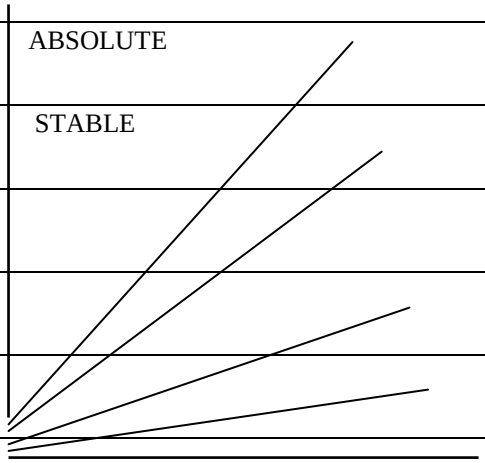
It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9>L_9-I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_9>L_9-I_9$, when $P_9^2/4>I$ (8a)

$P_9>L_9-I_9+\frac{(I_9-P_9^2}{(L_9-I)}$ when $P_9^2/4<I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Reffer couffesy Nunney

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability

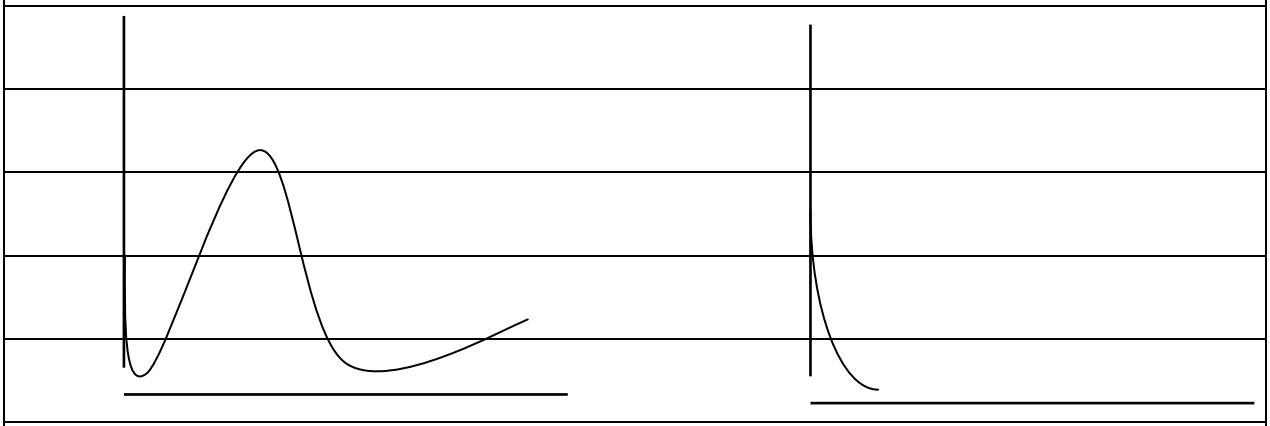
 

is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

Indeed, Nunny's model can be applied to some of the consciousness and philosophical problems. Predator prey analysis has ripped open the veneer of the fundamental concept of universe namely Debit-Credit world.

The noise of ordinary consciousness and (e&eb) beliefs drowns out (e) the signal. Unconsciousness is the background of (e) our ordinary awareness. Our organism is very much at (eb) the center of such effects. The organismic source is (=) our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is a meme-Scape. Pre-conceived concepts vie with structures, concepts with images. Like scientists who ignore assumed truths, we leapfrog over our beliefs and personality deficits, claiming idiosyncratic imagination is literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is both. We inherit the structures that make our experience what it is.

But the structure itself is “empty,” and each human culture “fills” it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us

feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through 're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns, nomenclature, and internal coherence of meaning and representation. In a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist

simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** The lower domain of experience is what we do not directly observe, the quantum realm, while the higher level is what we ordinarily observe, the classical realm. Based purely on experiment, the formulation of quantum theory initially placed the great divide between the quantum and classical domains between the measuring apparatus and the particle (Herbert, 1985). However, in the early 1930s, John von Neumann, in his rigorous mathematical treatise on quantum mechanics, found no support for such a division, compelling him to conclude that the wave for was collapsed by consciousness. Von Neumann's treatise has been called the "quantum bible" and "the most influential book on quantum theory ever written" (Herbert, 1985).

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t - T_9)F_9\big(R_9(t - T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t - T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ and the effective prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting $dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

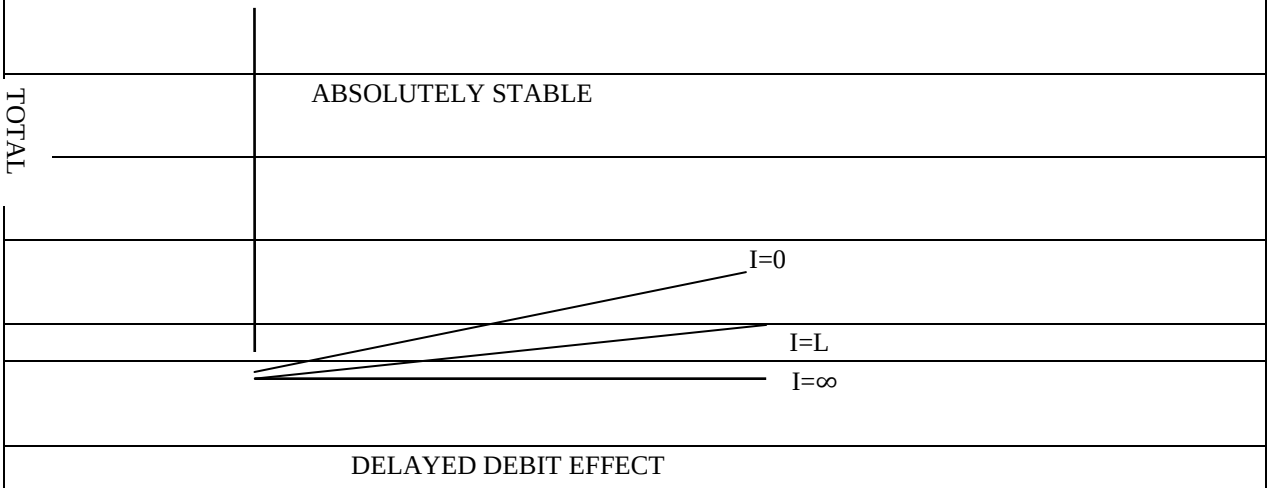



FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.

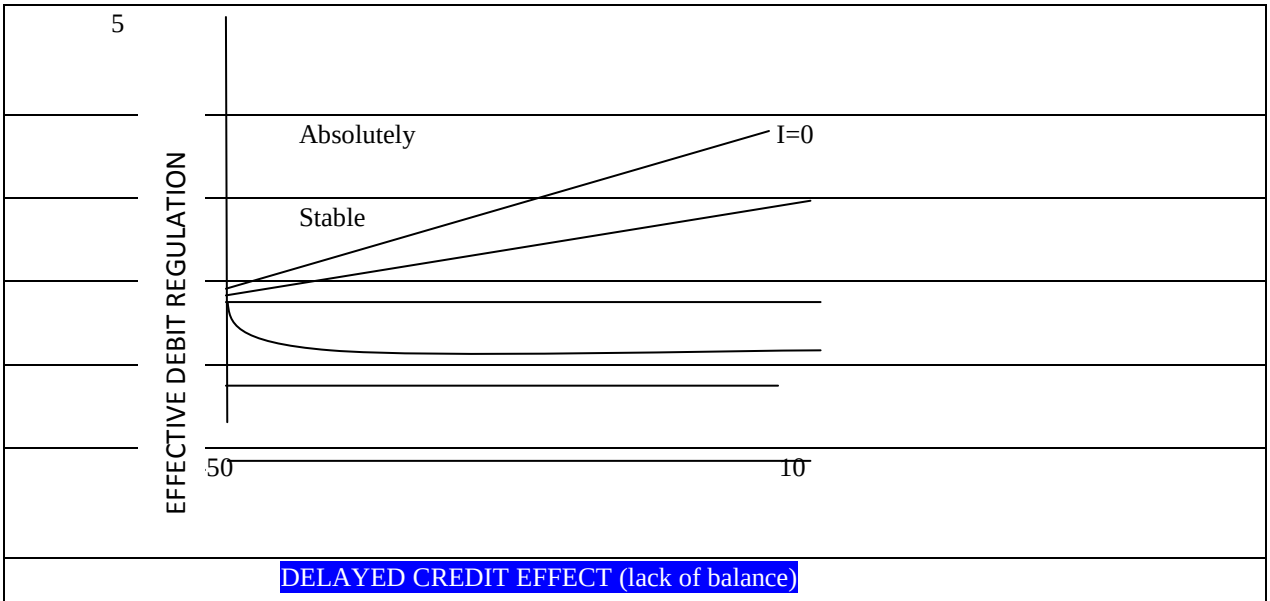



Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the

constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^\infty \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9\left(U_9 - V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9\right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$ then given the constraints of $f(T_9)$, $(A1.5)$ *can never be satisfied* Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDERATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 - V_9M_9 -]e^{-T_9\lambda_9} = 0$$

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9Z_9\sin(Z_9T_9) - (A_9M_9 + Y_9) + (A_9M_9 + X_9)\cos(Z_9T_9) = 0$$

$$(A_9M_9)Z_9 - M_9Z_9\cos(Z_9T_9) + (A_9M_9 - X_9)\sin(Z_9T_9) = 0$$

Where $A_9 = U_9 - V_9$ Using (A2.3), substitute for $Z_9^2$ and $Z_9$ in (A2.2):

$W_9^2(X_9^2 - 3M_9A_9X_9 - M_9^2(X_9 - Y_9) + 2M_9^3A_9 + 2A_9^2M_9^2]$

$$+W\left[MA^{3}-A^{2}\mathrm{X}+2\mathrm{MA}(2\mathrm{X}+\mathrm{Y})-2X^{2}\right]+A^{2}(\mathrm{X}+\mathrm{Y})=0$$

*where* $W_9$=1-cos ($\mathrm{Z}_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (\mathrm{X}_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ \mathrm{X}_9 - Y_9/2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range 0⍰$W_9$ ⍰2 *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (\mathrm{X}_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\ \mathrm{M}_9)^2\ (2A_9\mathrm{M}_9 - (\mathrm{X}_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9M_9^2 - 3M_9X_9\right) + W_9\left(3M_9A_9^2 - 2A_9X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9 + Y_9\right)$$

$=2M_9^2W_9^2(M_9 + A_9) + M_9A_9^2W_9 + \left(2A_9M_9 - X_9 + Y_9\right)(A_9W_9 + M_9W_9^2$

$+ (2- W_9)((2M_9X_9 + M_9Y_9)W_9 + A_9(X_9 + Y_9))$,

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range 0≤ $W_9$ ≤2.Defining

$Q_9(W_9, A_9) = a_9W_9^2 + c_9$ *where* $a_9$ $b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

**Harald Atmanspacher, Hans Primas** show that such an epistemic perspective can be consistently complemented by Einstein's ontically oriented position. **Time, Quantum and Information 2003, pp 301-321 Epistemic and Ontic Quantum Realities Harald Atmanspacher, Hans Primas Robert W. Spekkens** present a toy theory that is based on a simple principle: the number of questions about the physical state of a system that are answered must always be equal to the number that are unanswered in a state of maximal knowledge. Many quantum phenomena are found to have analogues within this toy theory. These include the noncommutativity of measurements, interference, the multiplicity of convex decompositions of a mixed state, the impossibility of discriminating nonorthogonal states, the impossibility of a universal state

inverter, the distinction between bipartite and tripartite entanglement, the monogamy of pure entanglement, no cloning, no broadcasting, remote steering, teleportation, entanglement swapping, dense coding, mutually unbiased bases, and many others. The diversity and quality of these analogies is taken as evidence for the view that quantum states are states of **incomplete knowledge** rather than states of reality. A consideration of the phenomena that the toy theory **fails to reproduce, notably, violations of Bell inequalities** and the existence of a Kochen-Specker theorem, provides clues for how to proceed with this research program.DOI: http://dx.doi.org/10.1103/PhysRevA.75.032110 **Evidence for the epistemic view of quantum states A toy theory Phys Rev A 75, 032110 – Published 19 March 2007 Robert W. Spekkens In such specialized and autonomous areas, it is to the eminent and erudite authors to interpret the contradistinctions and contradictions' in the dove tailing explaining. This I feel would be yeoman service in the creation of Nature's general ledger. M. S. Leifer and O. J. E. Maroney** examine the relationship between quantum contextuality (in both the standard Kochen-Specker sense and in the generalized sense proposed by Spekkens) and models of quantum theory in which the quantum state is maximally epistemic. We find that preparation noncontextual models must be maximally epistemic, and these in turn must be Kochen-Specker noncontextual. This implies that the Kochen-Specker theorem is sufficient to establish both the impossibility of maximally epistemic models and the impossibility of preparation noncontextual models. The implication from preparation noncontextual to maximally epistemic then also yields a proof of Bell's theorem from an Einstein-Podolsky-Rosen-like argument. DOI: http://dx.doi.org/10.1103/PhysRevLett.110.120401 Received 25 August 2012Published 20 March 2013 © 2013 American Physical Society **Maximally Epistemic Interpretations of the Quantum State and Contextuality Phys. Rev. Lett 110, 120401 – Published 20 March 2013 M. S. Leifer and O. J. E. Maroney.** Epistemological, onitic and ontological aspects particularly of quantum mechanics have invited

profound thinkers to contribute their mite to the growing literature on the subject, which includes the contributions from the subterranean realm and ceratoid dualism of philosophy like that Chaos theory of Deleuze, Derrida , Kierkegaard and others. **Andrei Khrennikov** shows that the Dirac-von Neumann formalism for quantum mechanics can be obtained as an approximation of classical statistical field theory. This approximation is based on the Taylor expansion (up to terms of the second order) of classical physical variables – maps f: Ω → R, where Ω is the infinite-dimensional Hilbert space. The space of classical statistical states consists of Gaussian measures ρ on Ω having zero mean value and dispersion σ2 (ρ) ≈ h. This viewpoint to the conventional quantum formalism gives the possibility to create generalized quantum formalisms based on expansions of classical physical variables in the Taylor series up to terms of nth order and considering statistical states ρ having dispersion σ2(ρ) = h n (for n = 2 we obtain the conventional quantum formalism). **December 2005, Volume 18, Issue 7, pp 637-650 Date: 28 Nov 2005 Generalizations of Quantum Mechanics Induced by Classical Statistical Field Theory Andrei Khrennikov** It is widely accepted that consciousness or, in other words, mental activity is in some way correlated to the behavior of the brain or, in other words, material brain activity. Since quantum theory is the most fundamental theory of matter that is currently available, it is a legitimate question to ask whether quantum theory can help us to understand consciousness. Several approaches answering this question affirmatively, proposed in recent decades, will be surveyed. It will be pointed out that they make different epistemological assumptions, refer to different neurophysiological levels of description, and adopt quantum theory in different ways. For each of the approaches discussed, these imply both problematic and promising features which will be indicated. **Discrete Dynamics in Nature and Society Volume 2004 (2004), Issue 1, Pages 51-73 http://dx.doi.org/10.1155/S102602260440106X Quantum theory and consciousness: an overview with selected examples Harald Atmanspacher Copyright © 2004 Hindawi Publishing**

**Corporation. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.** Another thought provoking and investigatory article is by Peter G. Lewis et al. Perhaps the quantum state represents information about reality, and not reality directly. Wave function collapse is then possibly no more mysterious than a Bayesian update of a probability distribution given new data. **Peter G. Lewis et al** consider models for quantum systems with measurement outcomes determined by an underlying physical state of the system but where several quantum states are consistent with a single underlying state—i.e., probability distributions for distinct quantum states overlap. Significantly, they demonstrate by example that additional assumptions are always necessary to rule out such a model. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.150404 **Distinct Quantum States Can Be Compatible with a Single State of Reality Phys. Rev. Lett. 109, 150404 – Published 9 October 2012 Peter G. Lewis et al** Communication complexity of a quantum channel is the minimal amount of classical communication required for classically simulating a process of state preparation, transmission through the channel and subsequent measurement. It establishes and perpetuates a limit on the power of quantum communication in terms of classical resources. **Alberto Montina** shows that classical simulations employing a finite amount of communication can be derived from a special class of hidden variable theories where quantum states represent statistical knowledge about the classical state and not an element of reality. This special class has attracted strong interest very recently. The communication cost of each derived simulation is given by the mutual information between the quantum state and the classical state of the parent hidden variable theory. Finally, find that the communication complexity for single qubits is smaller than 1.28 bits. The previous known upper bound was 1.85 bits. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.110501 **Epistemic View of Quantum States and**

**Communication Complexity of Quantum Channels Phys. Rev. Lett 109, 110501 – Published 12 September 2012 Alberto Montina for models see elsewhere .** Stability analysis of what? Quantum state representing reality or our knowledge of reality? Does the quantum state represent reality or our knowledge of reality? In making this distinction precise, **Nicholas Harrigan, Robert W. Spekkens** are led to a novel classification of hidden variable models of quantum theory. They show that representatives of each class can be found among existing constructions for two-dimensional Hilbert spaces. Approach also provides a fruitful new perspective on arguments for the nonlocality and incompleteness of quantum theory. Specifically, **Nicholas Harrigan, Robert W. Spekkens** show that for models wherein the quantum state has the status of something real, the failure of locality can be established through an argument considerably more straightforward than Bell's theorem. The historical significance of this result becomes evident when one recognizes that the same reasoning is present in Einstein's preferred argument for incompleteness, which dates back to 1935. This fact suggests that Einstein was seeking not just any completion of quantum theory, but one wherein quantum states are solely representative of our knowledge. Hypothesis is supported by an analysis of Einstein's attempts to clarify his views on quantum theory and the circumstance of his otherwise puzzling abandonment of an even simpler argument for incompleteness from 1927. **February 2010, Volume 40, Issue 2, pp 125-157 Date: 09 Jan 2010 Einstein, Incompleteness, and the Epistemic View of Quantum States Nicholas Harrigan, Robert W. Spekkens** statistical model of the probabilistic description of physical reality has been studied by many authors. **Andrei Khrennikov posits** a contextual statistical model of the probabilistic description of physical reality. Here contexts (complexes of physical conditions) are considered as basic elements of reality. There is discussed the relation with QM. He proposes a realistic analogue of Bohr's principle of complementarity. In the opposite of the Bohr's principle, **Andrei Khrennikov** principle has no direct relation with mutual exclusivity for

observables. To distinguish his principle from the Bohr's principle and to give better characterization, he changes the terminology and speaks about supplementarity, instead of complementarity. Supplementarity is based on the interference of probabilities. It has quantitative expression trough a coefficient which can be easily calculated from experimental statistical data. There is need not appeal to the Hilbert space formalism and noncommutativity of operators representing observables. Moreover, in our model there exists a pair of supplementary observables which cannot be represented in the complex Hilbert space. There are discussed applications of the principle of supplementarity outside quantum physics. **October 2005, Volume 35, Issue 10, pp 1655-1693 The Principle of Supplementarity: A Contextual Probabilistic Viewpoint to Complementarity, the Interference of Probabilities and Incompatibility of Variables in Quantum Mechanics Andrei Khrennikov Foundations of Physics** Description of the vacuum in Yang-Mills theory and arrive at a physical interpretation of the pseudoparticle solution and the attendant violation of symmetries has been done by **R. Jackiw and C. Rebbi.** It is here vaccuum stability forms the bastion, pillar, post and stylobate of the mechanisms of yang mills Theory and action of a particle. The existence of topologically inequivalent classical gauge fields gives rise to a family of quantum mechanical vacua, parametrized by a CP-nonconserving angle. The requirement of **vacuum stability** against gauge transformations renders the vacua chirally noninvariant. DOI: http://dx.doi.org/10.1103/PhysRevLett.37.172 **Vacuum Periodicity in a Yang-Mills Quantum Theory Phys. Rev. Lett 37, 172 – Published 19 July 1976 R. Jackiw and C. Rebbi.** Three-dimensional Yang-Mills and gravity theories augmented by gauge-invariant mass terms are analyzed. These topologically nontrivial additions profoundly alter the particle content of the models and lead to (eb) quantization of a dimensionless mass-coupling-constant ratio. The vector field excitations become (=) massive, with (e&eb) spin 1 (rather than massless with spin 0), and the mass provides (eb) an infrared cutoff. The gravitation

acquires (e) mass, mediates (e&eb) finite-range interactions, and has (e) spin 2 (rather than being absent altogether); although its mass term is of (e)third derivative order, there are no ghosts or acausalities. DOI: http://dx.doi.org/10.1103/PhysRevLett.48.975 © 1982 The American Physical Society **Three-Dimensional Massive Gauge Theories Phys. Rev. Lett 48, 975 – Published 12 April 1982 S. Deser et al** All scattering amplitudes in (eb) the maximally supersymmetric N=4 super-Yang–Mills theory possess (e) a new, dual superconformal symmetry which extends the previously discovered dual conformal symmetry of MHV amplitudes. To reveal this property authors formulate the scattering amplitudes as (=) functions on (e&eb) the appropriate dual superspace. Rewritten in this form, all tree-level MHV and next-to-MHV amplitudes exhibit manifest dual superconformal symmetry. Authors propose a new, compact and Lorentz covariant formula for (e) the tree-level NMHV amplitudes for (e) arbitrary numbers and types of external particles. The dual superconformal symmetry is broken at loop level by infrared divergences. However, authors provide evidence that the dual conformal anomaly of (e) the MHV and NMHV superamplitudes is (=) the same and, therefore, their ratio is (=) dual conformally invariant. Authors show this explicitly for (e) the six-particle amplitudes at one loop. Authors conjecture that these properties hold for all, MHV and non-MHV, superamplitudes in N=4 SYM both at (eb) weak and at strong coupling. **Nuclear Physics B Volume 828, Issues 1–2, 21 March 2010, Pages 317–374 Dual superconformal symmetry of scattering amplitudes in N=4 super-Yang–Mills theory J.M. Drummonda DOI: 10.1016/j.nuclphysb.2009.11.022.** By comparison with numerical results in (eb) the maximal Abelian projection of lattice Yang–Mills theory, it is argued that the nonperturbative dynamics of (e) Yang Mills theory can be described by (e) a set of fields that take their values in (eb) the coset space SU(2)/U(1). The Yang–Mills connection is parameterized in (eb) a special way to separate (e&eb) the dependence on the coset field. The coset field is then regarded as (=) a collective variable and a method to obtain its effective

action is developed. It is argued that the physical excitations of (e) the effective action may be (=) knot solitons. A procedure to calculate the mass scale of knot solitons is discussed for (e) lattice gauge theories in the maximal Abelian projection. The approach is extended to (e&eb) the SU (N) Yang–Mills theory. A relation between the large N limit and the monopole dominance is pointed out. **Physics Letters B Volume 458, Issues 2–3, 8 July 1999, Pages 322–330 An effective action for monopoles and knot solitons in Yang–Mills theory Sergei V. Shabanov DOI: 10.1016/S0370-2693(99)00612-7.** It is here the quintessential formulation of consciousness occurs. Like a man standing on the threshold of infinity trying to ponder what lies beyond the veil, which separates the seen from unseen a true seeker tries to lift the individual general ledger to the cosmic general ledger. Like in Darwinian evolution, mind also has to evolve from individual consciousness to cosmic consciousness. In Kaishmiri Shaivism one withdraws from the finite to the infinite, but one also goes on an outward journey from the infinite to the finite, because both the finite and the infinite have an intimate connection. The finite is not seen as unreal, but as a symbol of the infinite. There is no real distinction between them. Those two movements constitute Spanda, a key concept in Kaishmiri Shaivism. Spanda is the pulsation of the Absolute in different phases of being. There are no opposites like subject and object, unity or duality, absolute or relative. They are just different phases of the universal vibration of the Absolute. The goal is to realize or be (=) at once infinite and finite. One does not turn away from appearances (like in Advaita Vedanta), but one realizes that the Absolute manifests all things. Spanda, the eternal pulsation of the Absolute, oscillates between a passion to create and dispassion from the created. Through it the Absolute transform itself into all things and then returns back into the emptiness of its undifferentiated nature. In Kaishmiri Shaivism the Absolute is seen as pure consciousness (=being). The Absolute is an eternal all-pervasive principle, the highest reality, the nature of all entities eternally and blissfully at rest within its own nature. The Absolute is the nature of the Self (and

thus of us all). The Absolute is divine; it is Shiva, the Lord of the Universe. It is full of (e) conscious activity through (e&eb)which it generates the universe, and reabsorbs it into (e&eb) itself at the end of each cycle of creation. Thus we speak of monism, as everything resides within (eb) this one absolute consciousness. It sustains (eb) all things, it embraces (e) all things, and it pervades (e&eb) all things. All things are appearances within (eb) the absolute consciousness, but nevertheless real (in contrast to (e) Advaita Vedanta where (e) appearances are seen as unreal or illusionary). All things appear external (out there, outside ourselves), but they do not have (e) a being on their own. They do not exist as (=) separate entities on their own. Everything is contained within (eb) consciousness. What we see as objects are (=) manifestations of consciousness. The event which constitutes (eb) the universe is always internal events happening within (e&eb) consciousness because (e) their essential nature is consciousness itself. If a physical object were totally material, and independent or external to consciousness, it could never be experienced. The universe and (e&eb) consciousness are two aspects of a whole. The universe is (=) an attribute of consciousness which bears consciousness as (=) its substance. Doctrine of Vibration: S.G.Dyczkowski what you "see" is (=) not always what you "get." Many people mistakenly take their own visions literally without (e) "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in (eb) the physical sense but in a psychological sense. The archetypes of (e) the collective unconscious are arrayed behind (e) the scenes of current worldwide conditions, of (e) crisis and confusion. They mirror (e&eb) our own states back at us, whether we perceive them as (=) such or interpret them plausibly or not (Note the essay has been written in model able form and models are given in various papers.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4\,a_9 \square - b_9 \square 0$$

And

$$b_9^{\,2} \square 4\,a_9 c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9\, in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\, X_9$ ,

$Let\ Y_9 = \ y_9\, X_9$And try the substitution $A_9 = X_9/(2 + A_9\ y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + \ n_9 y_9)$ we have

$$W_9^{\,2}\, y_9 \left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^{\,2}\left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9\, X_9 \left[\frac{X_9 (1 + n_9 y_9)}{M_9^{\,2}(2 + n_9 y_9)^2} + 2\, y_9 \left(n_9 - 1\right) + \frac{X_9^{\,2}(9 + y_9)}{M_9^{\,2}(2 + n_9 y_9)} = 0.\ (A2.12)\right.$$

Since $n_9$ and $A_9$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9 = (2 + n_9 y_9)$, *the condition* (A2.99?) *can be written as*

$[\frac{X_9(E_9-9)}{E_9^2}]^2[\frac{X_9^2}{M_9^2} - 4\, y_9\ E_9^2 \frac{X_9}{M_9^2} + 4y_9 (n_9 - 2)E_9^3]$ .

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9 / W_9^2)$ can be expressed as

$2\, E_9 [y_9 E_9 + (y_9\ E_9 (9 + y_9)(E_9 - n_9))^{9/2}] > \frac{X_9}{M_9^2} > 2E_9 [y_9 E_9 - y_9 E_9 (9 + y_9)(E_9 - n_9))^{9/2}]$

Concept Of Being From Shaivite Point Of View: In His Famous Work Entitled On The Conservation Of Force (1847), Hermann Von Helmholtz Had Divided The Totality Of The Energy In The Universe Between Potential Energy And Kinetic Energy And Affirmed The Reciprocal Convertibility Of The Two. In 1852, William Thomson Pointed Out That There Exists A Sub-Ensemble Within Kinetic Energy, Heat, Which, Once It Has Been Generated, Is No Longer Entirely Convertible Into Potential Energy-Or Into Any Other Form Of Kinetic Energy. Considering That The (Partial) Reconversion Of Heat Into Labor Is Possible Only In Situations That Present A Disparity In Temperature, And That Heat Tends To Pass From Warmer To Cooler Bodies By Spreading On An Even Temperature Level Through Space, Thomson Concluded That The Universe Tends Towards A Final State Where Any Energetic Transformations, Every Movement And Every Form Of Life Will Cease: The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO: The Ultimate Nature Of Anything Cannot Change. Change Of Any Sort Is Merely Apparent (Vivarta). The World Of Change And Becoming Is A False Super- Imposition (Adhyaropa, Adhyasa) On The Absolute. It Appears So On The Screen Of Individual, Collective Consciousness, Albeit One Cannot Be So Sure Of Cosmic, For There Are Many Reasons That Cosmic Transactions Also Had To Be Subjected To Change

By ParaBrahaman. In Cosmic Terms, The Mistake (Bhranti) Consists Of The Supposition That The Real Brahman Is The Unreal Universe And The Unreal Universe Is The Real Brahman. In Microcosmic Terms, It Is The Mistake Of Falsely Conceiving The Body, Mind Or Even One's Personality To Be The Self. Note We Have Taken Self As Witness Consciousness. Notwithstanding There Is Nothing Wrong In The Conception Of Self As A Dynamical Presupposition, In That Individual Consciousness And Collective Consciousness And Even Cosmic Consciousness Evolutes In To ParaBrahaman. In The Same Way As The Image Of A Snake Is Falsely Superimposed On A Rope, Similarly The Universe Is Falsely Projected Onto The Real Substratum, The Brahman. Ignorance Is Not Merely A Personal Lack Of Knowledge, But A Cosmic Principle. As Such It Is Called "Maya," The Indefinable Factor (Anirvacaniya) That Brings This Mistake In Identity About. The Reality Status Of This Cosmic Illusion Is Also Indefinable: On The One Hand It Is Not Brahman, The Sole Reality; On The Other Hand It Is Not Absolutely Non-Existent Like A Hare's Horn Or The Son Of A Barren Woman. Note Here The Vedantin Avoids Talking About Zero Which We Have Avidly And Assiduously Stressed About. Nothing Is The Factor That Exists. Also See the Remarks about Satre's Being And Nothingness. Brahman Is The Source Of World Appearances Only In The Sense Of Being Their Unconditioned Ground Or Essential Nature. The Universe Is False Not Because It Has No Nature Of Its Own But Because It Does Have One. The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski. Concept Of Being From Shaivite Point Of View: It Is The 'I Am The Body/Mind' Belief That Gives Rise To The 'I Am Not The World' Belief. These Two Beliefs Are (=) Co-Created."Consciousness Projects The Appearance Of The Mind, Body And World By Taking (Eb) The Shape Of Thinking, Sensing And Perceiving."“Attention Is Consciousness With An Object. When The Object Vanishes, Attention Simply Remains What It Always Is Consciousness."“There Is No Purpose To Meditation. The Purpose Is Already Accomplished."“Everything

That Is Experienced Is Experienced By, Through, In And As Consciousness.”“The Seen Cannot Be Separated From Seeing And Seeing Cannot Be Separated From Consciousness.” “**The Reality Of Any Experience Is Not Hidden In The Appearance; It Is Expressed By The Appearance The Transparency Of Things**' By Rupert Spira Book Review By Dennis Waite The Following Is A Review Of Rupert's Book: 'The Transparency Of Things: Contemplating The Nature Of Experience’. There Is An Essay From The Book Here And The Essay On Cézanne May Be Read At The Awakened Eye Website Nature Rejects The Naiveté That Seeks Absolute Truth. We Are Beginning To Realize, Individually And Culturally, That "Realities" Are All Human Constructions. The Task Becomes One Of "Catching Ourselves In The Act" Of Creating Our Own "Reality" From The Flow Of Events. Human Truth Is Always An Engagement Of Mind With Experience. The Challenge Of The Therapist In These Times Of Chaotic Change Is To Validate The Concept That We Don't Have To Fear The Collapse Of What We Think We Are. Strange Attractors: Transference, Holography, And An **Archetype Burke, J. (2003). Strange Attractors: Transference, Holography, And An Archetype (Doctoral Dissertation, Pacifica Graduate Institute, 2003).** When we get to the topic and understanding of the Transcendental Logic, the brilliant and baffling Transcendental Deduction (TD) Recall the two movements just discussed, the one from experience to its conditions and the one from the forms of valid inference to the concepts that we must use in all judging (the Categories). This duality led Kant to his famous question of right (quid juris) (A84=B116): with what right do we apply the Categories, which are not acquired from experience, to the contents of experience? (A85=B117). Kant's problem here is not as arcane as it might seem. It reflects an important question: How is it that the world as we experience it conforms to our logic? In briefest form, Kant thought that the trick to showing how it is possible for the Categories to apply to experience is to show that it is necessary that they apply (A97).TD has two sides, though Kant never treats them separately John Marlado,

accounts, attributes, and ascribes the philosophy of Nishida KitarÅ is lucid, insightful, and deeply informative despite the highly questionable argument that gives the book its structure. In comparison with other, technically more accurate accounts of Nishida's philosophy, it is broader in scope than most but more single-minded in its approach. The analysis is for the most part limited to four books that have been translated into English, out of approximately 16 volumes of philosophical essays in Japanese, but it aims to capture the contours and most important details of Nishida's entire philosophy. To appreciate the achievement of this account, one would need only to read, in English translation or in the Japanese original, any of the volumes with their streams of repetitive, puzzling expressions and circuitous, often dead- end argumentation. Wilkinson's elucidation of details in the passages he samples makes it easier to understand why Nishida is considered Japan's greatest academic philosopher. A critical look at the frame of the book and some of its own claims reveals where, to my mind, it has gone astray. The argument that frames this book is repeated throughout: Nishida's philosophy is an attempt to articulate, in a characteristically Western philosophical manner, the Zen experience he had as a young man. Nishida first made the attempt by using the conceptual frameworks of Western philosophers, and after they failed him he developed his own, largely by recasting Buddhist insights. In the end, a philosophical articulation of the Zen worldview in Western terms or even in Eastern terms recast by conceptual, Western philosophical methods approved impossible. What Nishida attempted is incommensurable with Western philosophy. While the author first offers these statements as hypotheses (p. 2), one premise of his argument that there is such a thing as "the Zen world-view" or "the Zen experience “remains unquestioned; and another that Nishida's pivotal experience was the Zen experienceâis asserted as "beyond question" (p. 151; see also pp. 28, 48, 10, 114).Wilkinson's presentation of Zen in the first chapter is drawn largely from D. T. Suzuki, with help from Suzuki's portage Abe Masao; some comments on time, birth, and death in; and a few other sources. The

result is a caricature that imagines Zen as an invariable, unitary experience. In fact, Dagan's teachings on practice as the manifestation of enlightenment do not sit well with D. T. Suzuki's emphasis on a satori experience, and Suzuki does not enlighten us about what Zen in its various practices has historically been. The author would have benefited from acquaintance with other presentations by Zen teachers, such as that of ShunryÅ Suzuki'sZen Mind Beginner's Mind, with historical accounts of Zen traditions, and with contrasts between DÅgen's view of time and Sam Â· sÄraand that of other Buddhist philosophers. A little research in contemporary Zen scholarship would have shown how implausible is the talk of "nirvana or awareness of mu" (p. 24), "the basic Zen position [that] never varies" (p. 3), "the Zen vision" (pp. 5, 7), "Zen epistemology" , "Zen experience" (p. 59), "the conception of time involved in Zen" (p. 15), and "the fundamental Zen assertion that there is no absolute time" (p. 16). Even if DÅgen's notion that time and being are nondual (p. 16) epitomized most of Zen, Nishida's early notion of "a transcendent, unchanging reality apart from time" (p. 51) would contrast sharply with it. Equally questionable is Wilkinson's contention that Nishida's "pure experience" is equivalent to Zen satori. The notion of pure experience Nishida developed in his first major work may be pivotal for his entire philosophy, but he never claims or implies that any part of his philosophy presupposes an exceptional experience like satori that grounds his convictions and renders them unverifiable to those who have not had the experience, as Wilkinson intimates (pp. 55, 159, and 161). Nishida presents his notions of pure experience and intellectual intuition, the grasp of its unity, and later... Perennialism and Constructivism is studied by Randolph T Dible : experiencer or the observer and the transcendental subjectivity itself Pure objectivity is essentially a deconstruction, apophatic and negativa The encountering of the experiencer or observer—transcendental subjectivity itself—at the foundation of the world leads inevitably to the recognition of pure objectivity as ultimate reality (which can be taken as its ultimate deconstruction, analogous to the apophatic or via

negativa), from which objects derive their value, weight, significance, meaning or objectivity. In this way, pure objectivity can be seen as the supra-self-evident Axiological Axiom, so to speak, even Unconditional Love, in romantic terms. This axiology (value theory) has a structure inverse to the relationship between transcendental subjectivity as the radical unity of pure self-reference and on the other hand, the world of forms, as mere traces (representations, indications) of the unique, original "first distinction" Spencer-Brown speaks of at the foundation of his calculus. That is, all forms (i.e., distinctions, differences) would reduce to being the first distinction, also known as the marked state, which can be called penultimate reality (pure self-reference or transcendental subjectivity: the Spirit which animates us), except that forms are complimentary to their content, which is their objectivity or value, which would reduce to the unmarked state or ultimate reality. It is the incongruity of form (thoughts; Whitehead's "negative prehensions") and value (feelings; Whitehead's "positive prehensions," or my notion of objectivity, meaning and qualia; in short, the non-formal aspects of experience) that holds forms open and keeps them from absolute reduction. This accounts for the brute, concrete persistence of the "functional illusion"-- to use a term from Dzogchen Buddhism-- of the world. Thus this system has an axiology of metaphysical objectivity grounded on the ideal of pure objectivity as the source of all value, meaning and significance, itself the very fecundity of profundity, which is the motive of drawing the distinction in the first place **(Wikipedia, Kant's writings and Logic Of Sense)** Nano rings were created by accident while intending to make quantum dots. They have interesting optical properties associated with excitons and the Aharonov–Bohm effect. Application of these rings used as light capacitors or a buffer includes photonic computing and communications technology. Analysis and measurement of geometric phases in mesoscopic rings is ongoing. Several experiments, including some reported in 2012, show Aharonov-Bohm oscillations in charge density wave (CDW) current versus magnetic flux, of dominant period h/2e through CDW rings up to 85 **μm** in

circumference above 77 K. This behavior is similar to that of the superconducting quantum interference devices (seeSQUID). **Aharonov–Bohm nano rings (Wikipedia) Stability analysis of nano rings would help progressively accentuate the research about the materials division that uses nano rings.** From the profound contemplation and wisdom of the Buddha and Mahabodhisattvas, we know there is no such reality. Instead, egolessness (non-self) is the only path to understand the reality of the deluded life. All existences are subject to the law of causes and conditions. These include the smallest particles, the relationship between the particles, the planets, and the relationship between them, up to and including the whole universe! From the smallest particles to the biggest matter, there exists no absolute independent identity. Egolessness (non-self) implies the void characteristics of all existence. Egolessness (non-self) signifies (eb) the non-existence of permanent identity for self and existence (dharma). Sunyata stresses (eb) the voidness characteristic of (e) self and existence (dharma). Sunyata and Egolessness possess (e) similar attributes. As we have discussed before, we can observe the profound significance of Sunyata from (e) the perspective of inter-dependent relationships. Considering dharma-nature and (e&eb) the condition of nirvana, all existences are immaterial and of (e) a void-nature. Then we see each existence as (=) independent of each other. But then we cannot find any material that does exist independent of (e) everything else. So Egolessness also implies (eb) void-nature! From the observation of all existences, we can infer (eb) the theory of nirvana and the complete cessation of (e) all phenomena. From the viewpoint of phenomena, all existences are so different from (e) each other, that they may contradict (e&eb) each other. They are (=) so chaotic. In reality, their existence is (=) illusionary and arises from (e) conditional causation. They seem to exist on one hand, and yet do not (e) exist on the other. They seem to be united, but yet they are so different to (e) one another. They seem to exist and yet they do cease! Ultimately everything will return to (e) harmony and complete calmness. This is the nature of all existence. It is the

final resting place for all. If we can understand this reality and remove our illusions, we can find this state of harmony and complete calmness. **Teachings in Chinese Buddhism (6) sunyata (emptiness) in the Mahayana context (Wikipedia) Many models have been given concatenating Buddhist doctrines and Vaccum energy. These models essentially throw light on the essence of self(witness consciousness) and the cosmic general ledger (cosmic consciousness) while the imbalance in the gratification deprivation complex always remain unbalanced when it comes to individual general ledger , advocating the evolutionary theory of individual consciousness to cosmic consciousness. It is to be noted that it is not the question of knowing all the transactions but the wisdom of what lies behind the debit- Credit, Debit-Debit, or Credit-Credit principle that is most important for assimilation.** Look at this simple thesis to prove the point stated in the foregoing. Thesis: 'The World Has A Beginning In Time, And In Space It Is Also Enclosed In Boundaries.' Proof: 'For If One Assumes That The World Has No Beginning In Time, Then Up To Every Given Moment In Time An Eternity Is Elapsed, And Hence An Infinite Series Of States Of Things In The World, Each Following Another, Has Passed Away. But Now The Infinity Of A Series Consists Precisely In The Fact That It Can Never Be Completed Through A Successive Synthesis. Therefore An Infinitely Elapsed World-Series Is Impossible, So A Beginning Of The World Is A Necessary Condition Of Its Existence, Which Was The First Point To Be Proved (Kant, Critique Of Pure Reason, Tr. By R. Guyer And A. W. Wood (Cambridge: Cambridge University Press, 1998): B 454, P. 470). As Quoted In The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO"The Cyclical Hypothesis, So Heavily Criticized By Nietzsche (VP II 325 And 334), Arises In This Way." In Fact, Nietzsche Was Not Criticizing The Cyclical Hypothesis But Only The Particular Form Of That Hypothesis Presented In Vogt's Work. All Of Nietzsche's Texts Without Exception Speak Of The Eternal Return As The Repetitio Of The Same Events Within A Cycle Which Repeats Itself

Eternally. If Deleuze's Interpretation Holds That The Eternal Return Is Not A Circle, Then What Is It? A Wheel Moving Centrifugally, Operating A "Creative Selection," "Nietzsche's Secret Is That The Eternal Return Is Selective" Says Deleuze: The Eternal Return Produces Becoming-Active. It Is Sufficient To Relate The Will To Nothingness To The Eternal Return In Order To Realize That Reactive Forces Do Not Return. However Far They Go, However Deep The Becoming-Reactive Of Forces, Reactive Forces Will Not Return. The Small, Petty, Reactive Man Will Not Return. Affirmation Alone Returns, This That Can Be Affirmed Alone Returns, Joy Alone Returns. Everything That Can Be Denied, Everything That Is Negation, Is Expelled Due To The Very Movement Of The Eternal Return. We Were Entitled To Dread That The Combinations Of Nihilism And Reactivity Would Eternally Return Too. The Eternal Return Must Be Compared To A Wheel; Yet, The Movement Of The Wheel Is Endowed With Centrifugal Powers That Drive Away The Entire Negative. Because Being Imposes Itself On Becoming, It Expels From Itself Everything That Contradicts Affirmation, All Forms Of Nihilism And Reactivity: Bad Conscience, Ressentiment..., We Shall Witness Them Only Once. [...] The Eternal Return Is The Repetition, But The Repetition That Selects, The Repetition That Saves. Here Is The Marvelous Secret Of A Selective And Liberating Repetition. There Is No Need To Remind The Reader That Neither The Image Of A Centrifugal Movement Nor The Concept Of A Negativity-Rejecting Repetition Appears Anywhere In Nietzsche's Writings, And Indeed Deleuze Does Not Refer To Any Text In Support Of This Interpretation. Further, One Could Highlight That Nietzsche Never Formulates The Opposition Between Active And Reactive Forces, Which Constitutes The Broader Framework Of Deleuze's Interpretation. For Some Years, Marco Brusotti Has Called Attention To The Fact That Deleuze Introduced A Dualism That Does Not Exist In Nietzsche's Writings. To Be Sure, The German Philosopher Describes A Certain Number Of "Reactive" Phenomena (For Example, In The Second Essay Of The Genealogy Of Morality, § 11, He Talks About

"Reactive Affects" [Reaktive Affekte], "Reactive Feelings" [Reaktive Gefühlen], Reactive Men [Reaktive Menschen]); But These Are Nonetheless The Result Of Complex Ensembles Of Configurations Of Centers Of Forces That Remain In Themselves Active. Neither The Word Nor The Concept Of "Reactive Forces" Ever Appears In Nietzsche's Philosophy. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO Some Parts Are Deleted Due To Spatial Constraints. In the midst of thoughts about the eternal return we find at least two other thematic axes On the one hand, the view of the **world as a constant flux of forces without any goal, law, or rules of becoming. A chaos sive natura de-divinized and de-anthropomorphized which constitutes the "ontological substratum**" of the whole of Nietzsche's reflexions On the other hand, an ensemble of fragments of an anthropologico-sociological character, designing a path of liberation leading to the creation of superior individuals by way of a profound transformation of their instinctual structure. This transformation must be achieved by a practice of solitude and internal struggle towards the liberation from the ancient representations of the world and from the incorporated herd values. For an analysis of these thematic perspectives, see Paolo D'Iorio, La linea e il circolo. Cosmologia e filosofia dell'eterno ritorno in Nietzsche (Genova: Pantograf, 1995): 233-322. With most of the parameters quantified, with assumptions there must be picturesque view of the self and non self, space and time, science and philosophy with the concatenational consummations. After The Chapter "On Redemption," Where Zarathustra Dares Not Expose His Doctrine, The Eternal Return Begins To Be Enunciated In Part Three Of The Work. In The First Place, It Is The Dwarf Who Formulates It In The Chapter "On The Vision And The Riddle." Facing The "Gate Of The Instant" Which Symbolizes The Two Infinities That Stretch Towards The Past And The Future, The Dwarf Whispers: "All Truth Is Crooked, Time Itself Is A Circle" The Dwarf Represents The Spirit Of Gravity, And He Embodies The Herd Morality, "The Belittling Virtue" Which Is The Title Of Another Chapter From Part III. The Dwarf Can

Endure The Eternal Return Without Great Difficulties Because He Has No Aspirations; Unlike Zarathustra He Does Not Wish To Climb The Mountains That Symbolize Elevation And Solitude. In Two Unpublished Notes, From The Summer And The Fall Of 1883, Nietzsche Writes: The Doctrine Is At First Favored By The Rabble, Before It Gets To The Superior Men. The Doctrine Of Recurrence Will First Smile To The Rabble, Which Is Cold And Without Any Strong Internal Need. It Is The Most Ordinary Of Life Instincts, Which Gives Its Agreement First. Hence, The Content Of The Doctrine Is The Same, But Whereas The Dwarf Can Endure It (Because He Interprets It According To The Pessimistic Tradition For Which "Nothing Is New Under The Sun"), Zarathustra, Who Is The "Advocate Of Life" Regards The Eternal Return As The Strongest Objection To Existence, And As The Rest Of The Dream Suggests, He Does Not Yet Succeed In Accepting It After The Vision At The Gate Of The Instant, The Chapter Is Brought To An End By The Enigma Of The Shepherd. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO .How does the memory affect the GD complex or for that matter the gratification deprivation of an individual or a collective identity? Non-Markovian local-in-time master equations give a relatively simple way to describe the dynamics of open quantum systems with memory effects. Despite their simple form, there are still many misunderstandings related to the physical applicability and interpretation of these equations. Here, we clarify these issues both in the cases of quantum and classical master equations. Further introduction of the concept of a classical non-Markov chain signified through negative jump rates in the chain configuration. **E-M Laine et al 2012 J. Phys. B: At. Mol. Opt. Phys. 45 154004 doi:10.1088/0953-4075/45/15/154004 Local-in-time master equations with memory effects: applicability and interpretation.** Supposing That There Were Indeed An "Energy Of Contraction" Constant In All Centers Of Force Of The Universe, It Remains To Be Explained Where Any Difference Would Ever Originate. It Would Be Necessary For The Whole To Dissolve Into An Infinite Number Of

Perfectly Identical Existential Rings And Spheres, And We Would Therefore Behold Innumerable And Perfectly Identical Worlds COEXISTING [Nietzsche Underlines This Word Twice] Alongside Each Other. Is It Necessary For Me To Admit This? Is It Necessary To Posit An Eternal Coexistence On Top Of The Eternal Succession Of Identical Worlds? The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. Science is usually portrayed as dehumanizing. Brave New World epitomizes this fear. "The more we understand the world, the more it seems completely pointless" (Steven Weinberg). Certainly science can seem chilling when conceived in the abstract as a metaphysical world-picture. We may seem to find ourselves living in a universe with all the human meaning stripped out: Participants in a soulless dance of molecules, or harmonics of pointlessly waggling superstrings and their braneworld cousins. Nature seems loveless and indifferent to our lives. What rights have we to be happy? **(Brave New World: Aldous Huxley).** Religion can afford to be dogmatic but science cannot. Look at this passage which counters both scientists and philosophers alike: However, Hartmann Objects That The Regressive Movement Postulated By Schopenhauer Is Possible Only In Thought: It Remains Nothing More Than An "Ideal Postulate" With No Real Object And Which "Does Not Teach Us Anything About The Real Process Of The World That Unfolds In A Movement Contrary To This Backwards Movement Of Thought" (Hartmann, Philosophie Des Unbewussten, Third Edition (1871): 772). Hartmann Affirms That If Unlike Schopenhauer One Admits The Reality Of Time And Of The World Process, One Must Also Admit That The Process Must Be Limited In The Past And Therefore That There Must Be An Absolute Beginning. In Hartmann's Mind, Failure To Do So Would Result In Positing The Contradictory Concept Of An Accomplished Infinity: "The Infinity That From The Point Of View Of Regressive Thinking, Remains An Ideal Postulate, Which No Reality May Correspond To, Must, For The World, Whose Process Is, On The Contrary, A Progressive Movement, Open Up To A Determinate Result; And Here The Contradiction Comes To Light" (Hartmann

(1871): 772). What Really "Comes To Light" In This Passage Is The Fact That Hartmann Does Not Provide A Demonstration But A Petitio Principii. Indeed, The Concept Of The World Process Analytically Contains The Concept Of A Beginning Of The World. In All Rigors, It Is Therefore Impossible To Demonstrate These Concepts With Reference To Each Other. Secondly, Hartmann's View That One Is Bound To Accept The Reality Of The World Process Even If One Rejects The Ideality Of Schopenhauer's Time Is Mistaken. Hartmann Believes That If Time Is Real There Must Be A World Process With Both An Absolute Beginning And An Absolute End. Without Any Justification, Hartmann Jumps From Schopenhauer's Negated Time To Oriented Time. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. In the last years several theoretical papers discussed if time can be an emergent property deriving from ) quantum correlations. Here, to provide an insight into how this phenomenon can occur, authors present an experiment that illustrates Page and Wootters' mechanism of "static" time, and Gambini et al. subsequent refinements. A static, entangled state between a clock system and the rest of the universe is perceived as evolving by internal observers that test the correlations between the two subsystems. They implement this mechanism using an entangled state of the polarization of two photons, one of which is used as a clock to gauge the evolution of the second: an "internal" observer that becomes correlated with the clock photon sees the other system evolve, while an "external" observer that only observes global properties of the two photons can prove it is static.Phys. Rev. A 89, 052122 (2014)DOI:10.1103/PhysRevA.89.052122 arXiv: 1310.4691 [quant-ph] **Time from quantum entanglement: an experimental illustration Ekaterina Moreva.** Time is an entanglement .When the new ideas of quantum mechanics spread through science like wildfire in the first half of the 20th century, one of the first things physicists did was to apply them to gravity and general relativity. The results were not pretty. It immediately became clear that these two foundations of modern physics were entirely

incompatible. When physicists attempted to meld the approaches, the resulting equations were bedeviled with infinities making it impossible to make sense of the results. Then in the mid-1960s, there was a breakthrough. The physicists John Wheeler and Bryce DeWitt successfully combined the previously incompatible ideas in a key result that has since become known as the Wheeler-DeWitt equation. This is important because it avoids the troublesome infinites—a huge advance. But it didn't take physicists long to realise that while the Wheeler-DeWitt equation solved one significant problem, it introduced another. The new problem was that time played no role in this equation. In effect, it says that nothing ever happens in the universe, a prediction that is clearly at odds with the observational evidence. This conundrum, which physicists call 'the problem of time', has proved to be thorn in flesh of modern physicists, who have tried to ignore it but with little success. Then in 1983, the theorists Don Page and William Wootters came up with a novel solution based on the quantum phenomenon of entanglement. This is the exotic property in which two quantum particles share the same existence, even though they are physically separated. Entanglement is a deep and powerful link and Page and Wootters showed how it can be used (e) to measure time. Their idea was that the way pair of entangled particles evolves is a kind of clock that can be used to measure change. But the results depend on how the observation is made. One way to do this is to compare the change in the entangled particles with an external clock that is entirely independent of the universe. This is equivalent to god-like observer outside the universe measuring the evolution of the particles using an external clock. In this case, Page and Wootters showed that the particles would appear entirely unchanging—that time would not exist in this scenario. But there is another way to do it that gives a different result. This is for an observer inside the universe to compare the evolution of the particles with the rest of the universe. In this case, the internal observer would see a change and this difference in the evolution of entangled particles compared with everything else is an important a measure of time. This is

an elegant and powerful idea. It suggests that time is an emergent phenomenon that comes about because of the nature of entanglement. And it exists only for observers inside the universe. **Time is an entanglement** (Wikipedia). Danish religious philosopher Soren Kierkegaard carried out a systematic critique of the pretensions of reason and an abstract rationalism which he believed that the modern age was nurturing. Condemning reflection as a "danger" that ensnares people in logical delays and machinations, Kierkegaard compared it to a prison. Reflection is for him a form of captivity, a bondage which "can only be broken by [passionate] religious inwardness" (1978: 81). **Reflection seduces individuals into thinking its possibilities are "much more magnificent than a paltry decision"** (1978: 82). It leads them to act **"on principle,"** to dwell on the deliberation of the context of their actions and the calculation of their worth or outcome. Kierkegaard argues that this drives away feeling, inspiration, and spontaneity, all of which are crucial for true inner being and a vital relation to God. For Kierkegaard, as Nietzsche would later agree, genuine inner being (and culture) is characterized by the tautness and tension of the soul which characterizes passionate existence. But the "coiled springs of life relationships ... lose their resilience" in reflection (1978: 78) and "everything becomes meaningless externality, devoid of [internal] character" (1978: 62). Kierkegaard thus contributes to the development of an irrationalist tradition that has echoes in some later postmodern thought. Kierkegaard might well have agreed with his contemporary Fyodor Dostoyevsky, who wrote: "An intelligent [reflective] man cannot seriously become anything ... excessive consciousness is a disease" (1974: 3, 5). In an age overtaken by rules and regulations, genuine action -- which Kierkegaard assumes to be subjective and spontaneous -- is frustrated at every turn. Complaining that we are too "sober and serious" (1978: 71) even at banquets, Kierkegaard bemoans the fact that even suicides are premeditated (1978: 68)! "That a person stands or falls on his actions is becoming obsolete; instead, everybody sits around and does a brilliant job of bungling through with the aid of some reflection

and also by declaring that they all know very well what has to be done" (1978: 73). Thus, it is passion, not reflection, that guarantees "a decent modesty between man and man [and] prevents crude aggressiveness" (1978: 62). "Take away the passion and the propriety also disappears" (1978: 64). The ambiguity in the word "passion" may cause some confusion here. To say that the age and its individuals are "passionless" is not to say there are no emotions whatsoever, but rather that there is no true spiritual inwardness and depth, no intensively motivated action and commitment. It suggests that passion exists only in a simulated, pseudo-form, "the rebirth of passion" through "talkativeness" (1978: 64). "Chattering" for Kierkegaard gets in the way of "essential speaking" and merely "reflects" inconsequential events (1978: 89-99). Hence, in "the present age," emotions -- which in fact are all too pronounced -- have been transformed into negative forces. Anticipating Nietzsche's genealogy of the "slave revolt" in morality, Kierkegaard claims that the "enthusiasm" of the prior age of Revolution, a "positively unifying principle," has become a vicious "envy," a "negatively unifying principle" (1978: 81), a leveling force in its own right insofar as those lacking in talent and resources want to tear down those who have them. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS By Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html.** In November 2011, the United States government argued before the US Supreme Court that it wants to continue utilizing GPS tracking of individuals without first seeking a warrant. In response, Justice Stephen Breyer questioned what this means for a democratic society by referencing Nineteen Eighty-Four. Justice Breyer asked, "If you win this case, then there is nothing to prevent the police or the government from monitoring 24 hours a day the public movement of every citizen of the United States. So if you win, you suddenly produce what sounds like 1984... "[65]. What is to be said in unmistakable terms is that when it comes to national security and innocent citizens there is question of compromise even it be an infringement and encroachment of so called

privacy which has been made a word to subterfuge the activities that are far said to be fair. Quantum theory has provoked intense discussions about its interpretation since its pioneer days. One of the few scientists who have been continuously engaged in this development from both physical and philosophical perspectives is Carl Friedrich von Weizsäcker. The questions he posed were and are inspiring for many, including the authors of this contribution. Weizsäcker developed Bohr's view of quantum theory as a theory of knowledge

## ABSOLUTE STABILITY

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_9$ ) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9 - 9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2 - 2y_9\frac{(E_9 + n_9^2 y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right) - 4y_9(E_9 - n_9)$$

And has the solution

$$\frac{X_9}{M_9^2} > \frac{E_9}{(E_9 - 9)}[y_9(E_9 + n_9^2 y_9) + [y_9^2 (E_9 + n_9^2 y_9)^2 + 4y_9(E_9 - 9)(E_9 - n_9)]^{9/2}]$$

Given $E_9 > n_9$ this lower limit to $X_9 / M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9 / M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$ in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9 < n_9$ $(ie., 2A_9 M_9^2 < X_9 - Y_9)$ then (A2.9)and (A2.8)show that $Q_9$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2} > 2E_9[y_9 E_9 - (y_9 E_9 (9 + y_9)(E_9 - n_9)]^{9/2}]$$

And

$$E_9 > n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2 /2M_9 (X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$ and $X_9$ and $Y_9$ by $W_9^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9$ *and* separate (A3.1)into real and imaginary parts

$$Z_9^2 + (M_9 + V_9) Z_9 \sin(Z_9 T_9) - (U_9 M_9 + Y_9) + (M_9(U_9 + V_9) - X_9)$$

$$\mathrm{X}\cos(Z_9 T_9) - 2M_9 V_9 \cos^2(Z_9 T_9) + M_9 V_9 = 0$$

$$(M_9 + U_9) Z_9 - (M_9 + V_9) Z_9 \cos(Z_9 T_9) + (M_9(U_9 + V_9) - X)$$

$$\mathrm{X}\sin((Z_9 T_9) - M_9 V_9 \sin(Z_9 T_9) = 0$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 = 1 - \cos Z_9 T_9$and $A_9 = U_9 - V_9$we obtain

$$4 M_9^2 V_9^2 W_9^4 + 2 M_9 V_9 W_9^3[(M_9 - V_9)^2 + A_9(M_9 + V_9) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2[M_9 - V_9)^2(2 A_9 M_9 + Y_9 - X_9) + (2A_9 M_9 - X_9)^2 - A_9(M_9 - V_9)(2A_9 M_9 - X_9)$$

$$+ M_9 V_9 (A_9^2 + 4(X_9 + Y_9))]$$

$$+ W_9[A_9^2 + 2X_9)(2 A_9 M_9 - X_9) + 2A_9 Y_9(M_9 + V_9) - M_9 A_9^3] + A_9^2(X_9 M_9 - Y_9) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9\left(2, U_9 V_9\right) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting cos $(Z_9 T_9) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MDDULE NUMBERED SEVEN**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability

criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

## Morphological And Epistemological And Syllogistic Models For Which Nunny's Model Could Be Applied And Has Been Applied:

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production**

**Charge Asymmetry, $\Delta Y \equiv Yt - Yt^{-}$, Reflects The Asymmetry In $Tt^{-}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles

ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic

quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation

between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to

nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of

Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first

impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of

another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease. Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation

of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces

always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the

thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of

thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity

and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI

Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs” Outcomes—

And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities

Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \quad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

## DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t-T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\big(R(t)\big)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and haimovici model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

**TIME DELAY**

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Limitations of Predator prey analysis in immense. When we realize that world itself is Debit-Credit world and sometime credit-Credit world (See an earlier paper entitled the same nomeneclature.

Opposite limit of absolute stability. The result is a specific The Doctrine of Vibration awareness of an object of the form: 'here this is such.' While individual appearances do not lose their separate identity {svarupabheda) when they rest on a common basis, even so the particular object which appears according to its own characteristics (svalak$ana) is an individual reality in its own right. It is a different kind of appearance characterised by its association with the appearance of the specific location and time in which it is made manifest. The form of our experience is thus 'I now see this here'. But when we perceive each particular constituent appearance separately, each assumes a separate fixed function. Abhinava cites the following colourful example to illustrate how the various combinations of appearances account for the variety of experience: Thus even though the appearance of the beloved may manifest externally, it is as if far away in the absence of another appearance, namely, that of 'embracing'. So when the [appearing of the beloved] is associated with another appearance [namely that of 'far away'] the power (arthakriya) it formerly had of giving pleasure appears as its contrary. The form our experience assumes depends, not only on the nature of the object perceived, but also on personal factors entirely peculiar to ourselves. This theory explains this in two ways. In one sense, the object remains the same, but one or other of its constituent appearances

comes to the fore according to the inclinations of the perceiver. From another point of view, we can say that the perceived object is different for each perceiver according to the difference in the prominent appearance manifest to him. Abhinava, citing as an example a golden jar, illustrates how the same object appears differently to different perceivers according to the use they wish to make of it and to their state of mind: When a person who is depressed and feels that there is nothing [of value for him in the world] sees the jar, he merely perceives the appearance 'exists' [in the form of the awareness that] 'it is\ He is not conscious of any other [of its constituent appearances] at all. An individual who desires to fetch water [perceives] the appearance 'jar'. The man who simply wants something that can be taken somewhere and then brought back [perceives] the appearance 'thing'. The man who desires money [perceives] the appearance 'gold'. The man who desires a pleasing object [perceives] the appearance 'brightness' while he who wants something solid sees the appearance 'hardness'. These 'atomic events' or appearances emerge from the pure subject's consciousness and combine together to form a total event at each moment. Integral Monism of Kashmiri Saivism Daily life (vyavahara) goes on by virtue of this ever renewed flux of appearances. They are connected together and work towards a single unified experience because they appear within the field of consciousness of the universal subject. The aggregate of appearances arises in the [supreme] subject as do [sprouts in] a rice field. Even though each sprout germinates from its own seed, they are perceived as a collective whole. Appearances rest in this way within the universal subject. 'External- ity' is itself another appearance; it arises from a distinction between appearances and the individual subject. So, although all manifestation always occurs within the subject, it appears to be external due to the power of Maya which separates the individual subject from his object. This split must occur for daily life to be possible. Only externally manifest appearances can perform their functions; when they are merged within the subject and at one with him, they cannot do so. Daily life proceeds on the basis of the operation and withdrawal of the conditions necessary for fruitful action to be possible. Appearance in this sense represents the actualisation of a potential hidden in consciousness made possible by virtue of its dynamic, Spanda

nature which is both the flow from inner to outer and back as well as the power that impels it. The emergence from, and submergence into, pure consciousness of each individual appearance is a particular pulsation (visesaspanda) of differentiated awareness. Together these individual pulsations constitute the universal pulse (samanyaspandd) of cosmic creation and destruction. Thus, every single thing in this way forms a part of the radiant vibration {sphuratta, sphurana) of the light of absolute consciousness. Light and Awareness: The Two Aspects of Consciousness Absolute consciousness understood as the unchanging ontological ground of all appearing is termed 'Prakasa'. As the creative awareness of its own Being, the absolute is called ' Vimarsa '. Prakasa and Vimarsa — the Divine Light of consciousness and the reflective awareness this Light has of its own nature — together constitute the all-embracing fullness (purnata) of consciousness. The Recognition (pratyabhijfia) school of Kashmiri Saivism develops this concept of the absolute which finds its fullest expression in Utpaladeva's Stanzas on the Recognition of God. Even though neither of these two key terms appear in the Stanzas on Vibration or the Aphorisms of Siva, they recur frequently in their commentaries. Thus, although the original formulation of the Doctrine of Vibration differs from the theology of Recognition in this respect, it was extended in the course of its development to accommodate this concept of the absolute as well. This was possible, and quite justified, insofar as the absolute understood in Pratyabhijfia terms does not, as we shall see, differ essentially from that of the Spanda school. We can, as Kashmiri Saivites themselves have done, explain one in terms of the other. The Doctrine of Vibration Prakasa: The Light of Consciousness Prakasa is the pure luminosity' (bhdna) or 'self-showing' that constitutes the essence and ultimate identity (atman) of phenomena. That things appear at all is due to the light of consciousness, and their appearing (avabhasana) is itself this Light which bestows on all things their evident, manifest nature. Established in the light of consciousness everything appears there according to its own specific nature (svabhava). Anything that supposedly does not rest in this Light is as unreal as a sky-flower. 3 Thus, according to Rajanaka Rama, unlike the light of the sun, or any other light, this Light not only makes all things apparent, it is also their ultimate source. 4 Full of its

divine vibration the Light makes all things manifest and withdraws them into itself. This supra- temporal activity characterises it most specifically; devoid of it, it would be no better than an inert physical phenomenon. At the same time, this light is the conjunction (slesa) or oneness (aikdtmya) of its countless manifest forms, 6 and the collective whole (Sampina'ana) of all the categories of existence. The universe is nothing but the shining of the Light within itself. It is the radiant vibration (sphuratta) of this Light, the state (avast hand) in which consciousness becomes manifest. Although the Light shines as all things at all times and hence also makes their diversity manifest, 9 penetrating each object individually as well as collectively, it is not totally 'merged' (magna) or identified with the object so as to suffer any division within itself. Our experience of any object is of the form: T see this': it is not itself an object, but the manifest form the object assumes as a luminous principle of experience. The Light is ever revealed and can never be obscured; objectivity can never cast a shadow on the light of consciousness. The Stanzas on Vibration declare: That in which all this creation is established and from whence it arises is nowhere obstructed because it is unconditioned by [its very] nature. This Light is the highest reality (Paramartha). lity with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**DELAY DUE TO CONFEDERATION DEBIT TYPE MATURATION CREDIT TYPE CONSOLIDATION MATURATION**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \quad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9} + \frac{N_9 dG_9}{dR_9} - \frac{dB_9}{dR_9} > \frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9} + \frac{dM_9}{dR_9}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened**

**eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production Charge Asymmetry, $\Delta Y \equiv Yt - Y\bar{t}$, Reflects The Asymmetry In $T\bar{t}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton

particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results

in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely

dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact**

**most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the**

**thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is

nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the

world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease.

Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition

and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the

second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular

chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon

be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining

Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been

Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing

Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

**ABSOLUTE STABILITY**

$$L_9 = \frac{N_9 G_9}{2M_9} \frac{d(9nF_9)}{dR_9} (\geq 0) \quad (9a)$$

$$I_9 = -\frac{N_9 G_9}{2M_9} \frac{d(9nM_9}{dR_9} (\geq 0) \quad (9b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9 > L_9 - I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_9 > L_9 - I_9$, when $P_9^2/4 > I$ (8a)

$P_9 > L_9 - I_9 + \frac{(I_9 - P_9^2}{(L_9 - I)}$ when $P_9^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to

reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

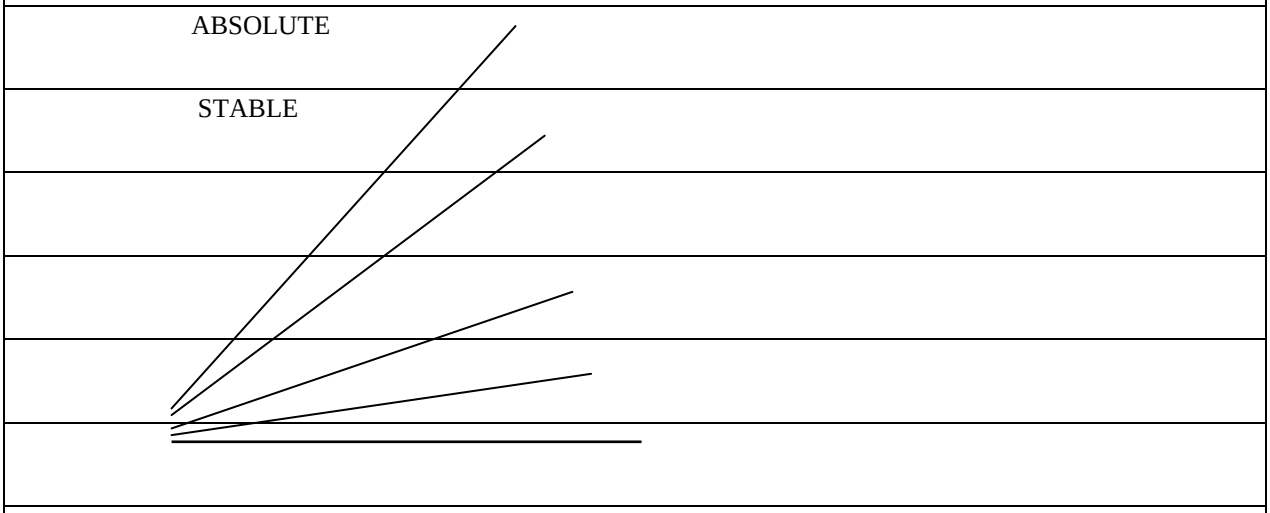



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori

reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

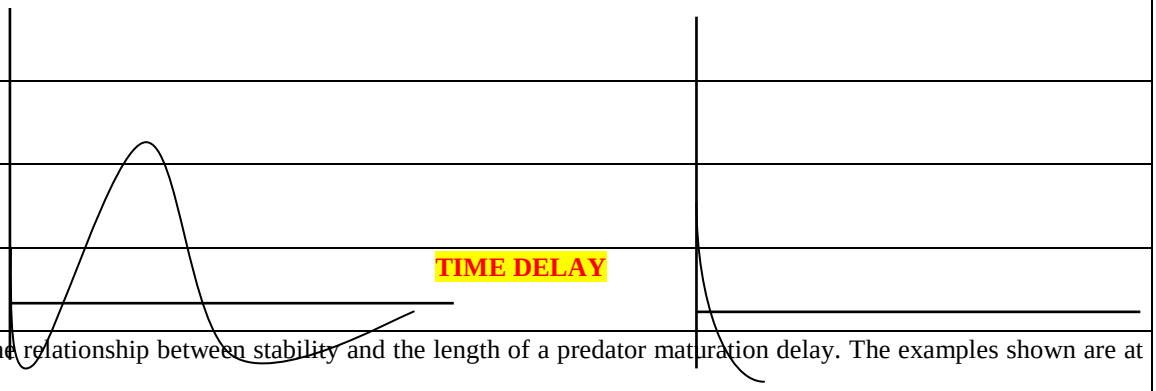



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

**DELAY IN BOTH PREDATOR AND PREY VIS A VIS DEBIT TYPES AND CREDIT TYPES**

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t-T_9)F_9\big(R_9(t-T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t-T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) and the effective prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting

$dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

Systems to which model has been applied also include: **Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI:

http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors.

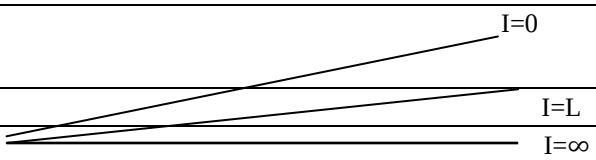
Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, ,**

**Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)**

TOTAL

ABSOLUTELY STABLE

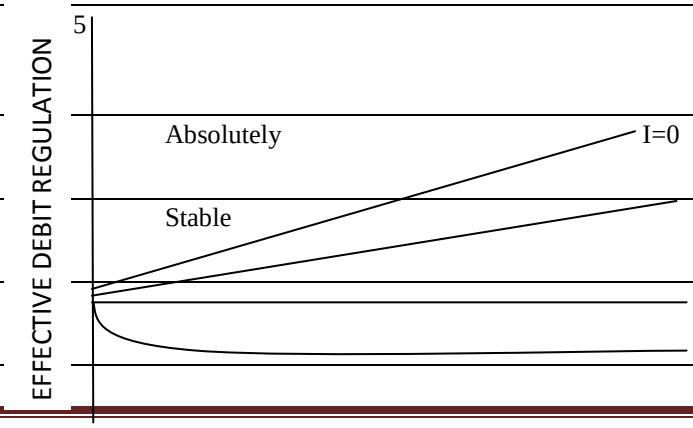



DELAYED DEBIT EFFECT

FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.



-50 10

**DELAYED CREDIT EFFECT (LACK OF BALANCE) CREDIT NON CUMSUMMATION:"**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

The following treatise bears ample testimony and infallible observatory to the fact that how Nunny's model can be applied to various cosmological, general relativity and economic problems.

## ONTOLOGICAL AND EPISTEMOLOGICAL FEATURES OF THE STUDY: OF NUNNY's MODEL

When they are known to be as they are, in the fullest sense of their existence, their phenomenal nature disappears leaving the ground of Being naked and accessible. This approach was validated by a critique of experience. The Vedanta established that space, time and the other primary categories of our daily experience can have no absolute existence. It was therefore necessary to make a distinction between relative truths — that accepted by the precritical common man — and an absolute truth discovered at a higher level of consciousness. The Saiva absolutist 17 rejects any theory that maintains that the universe is less than real. From his point of view a doctrine of two truths, one absolute and the other relative, endangers the very foundation of monism. The Kashmiri Saiva approach is integral: everything is given a place in the economy of the whole. It is equally wrong to say that reality is either one or diverse. Those who do so fail to grasp the true nature of things which is neither as well as both. "We do not" says Abhinavagupta, "base our contention that [reality] is one because of the contradictions inherent in saying that it is dual. It is your approach {paksa) that accepts this [method] While, if [duality and oneness] were in fact [to contradict each other], they would clearly be two [distinct realities]." The Vedantin, who maintains that non-duality is the true nature of the absolute by rejecting duality as only provisionally real, is ultimately landed in a dualism between the real and illusory by the foolishness of his own excessive sophistry (vacafadurvidya). Oneness is better understood as the coextensive unity (ekarasa) of both duality and unity. They are equally expressions of the

absolute. Gopinath Kaviraj says: According to Sarikara, Brahman is truth and Maya is inexplicable (anirvacaniya). Hence the [Advaitin's] endeavour to demonstrate the superiority of Advaita philosophy is turned against his own system. It tarnishes the picture of its philosophical perfection and profundity. He cannot accept Maya to be a reality, therefore his non-dualism is exclusive. The whole system is based on renunciation and elimination and thus is not all-embracing.... By accepting Maya to be Brahman (brahmamayi), eternal (nitya) and real (satyarupa), Brahman and Maya [in the Tantra] become one and coextensive. **The Doctrine of Vibration: An Analysis of the Doctrines and Practices of Kashmir Shaivism Mark S. G. Dyczkowski. Concept of being from Shaivite point of view:** Solutional behaviour and stability analysis of such systems is of paramount importance, notwithstanding the assumptions made which are not farfetched, in that it throws light on the fact whether they could be candidates for yet unsolved problems, providing Authetic ative determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication for the wide amplitudinal ramificatory usage of the thesis propounded. The world is now the other person's world, a foreign world that no longer comes from the self, but from the other. The other person is a "threat to the order and arrangement of your whole world...Your world is suddenly haunted by the Other's values, over which you have no control"(Satre: being and Nothingness) Sartre states that many relationships are created by people's attraction not to another person, but rather how that person makes them feel about themselves by how they look at them. This is a state of emotional alienation whereby a person avoids experiencing **their subjectivity by identifying themselves with "the look" of the other**. The consequence is **conflict**. In order to maintain the person's own being, the person must control the other, but must also control the **freedom of the other "as freedom**". These relationships are a profound manifestation of "bad faith" as for-itself is replaced with the other's freedom. The purpose of either participant is not to exist, but to maintain the other participant's looking at them. This system is often mistakenly called "love", but it is, in fact, nothing more than emotional alienation and denial of freedom. In

essence we are always simulating and selling ourselves at the cost of our freedom. Both as a script writer and participant consciousness, human beings experienceâ complete subjective feelings and sometimes absolute subjective feelings, which if remains unchecked might lead to pathological ramifications. (Italics mine) Perspicuous forbearance, Sophisticated seasoning, Partcipational observation, Background actuality, Existential worldliness, Existential strife, Predicational anteriorities, Character consonance, Ontological consonance, Primordial exactitude, Phenomenological correlates, Accolytish representation, Atrophied asseveration. Anamensial alienisms, Anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orinetatilities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state form the bastion, pillar, post, stylobate and sentinel of the classification scheme and doxa. appropriate and diagnostic, differentiating-discriminating, distinctive-distinguishing, emblematic- especial, essentially exclusive, fixed, idiosyncratic-inborn, inbred, indicative and individual, individualistic-inherent, individualizing and ingrained, inherent, innate, local, marked, native, normal, original, particularly- peculiar, personal, private, proper, regular, representative-singular, special, specific, symbolic, symptomatic- unique these form the classification and stratification measure that is suitable notwithstanding the constancy in many a case, as in gravity, which depends on the mass and distance for two given objects**.** Like Kant and Bergson, Deleuze considers traditional notions of space and time as unifying forms imposed by the subject. He therefore concludes that pure difference is non-spatio-temporal; it is an idea, what Deleuze calls "the virtual". (The coinage refers to Proust's definition of what is constant in both the past and the present: "real without being actual, ideal without being abstract.") Proust, Le Temps Retrouvé, ch. III: see the fourth line from the bottom of this page, or, in English translation, the thirteenth paragraph here: "I began to discover the cause by comparing those varying happy impressions which had the common quality of being felt simultaneously at the actual moment and at a distance in time, because of which common quality the noise of the spoon upon the plate, the unevenness of the paving-stones, the taste of the madeleine, imposed the past

upon the present and made me hesitate as to which time I was existing in. Of a truth, the being within me which sensed this impression, sensed what it had in common in former days and now, sensed its extra-temporal character, a being which only appeared when through the medium of the identity of present and past, it found itself in the only setting in which it could exist and enjoy the essence of things, that is, outside Time. [...] **Nothing but a moment of the past Much more perhaps; something which being common to the past and the present, is more essential than both. [...] a marvellous expedient of nature had caused a sensation to flash to me—sound of a spoon and of a hammer, uneven paving-stones**—simultaneously in the past which permitted my imagination to grasp it and in the present in which the shock to my senses caused by the noise had effected a contact between the dreams of the imagination and that of which they are habitually deprived, namely, the idea of existence—and thanks to that stratagem had permitted that being within me to secure, to isolate and to render static for the duration of a lightning flash that which it can never wholly grasp, a fraction of Time in its pure essence. When, with such a shudder of happiness, I heard the sound common, at once, to the spoon touching the plate, to the hammer striking the wheel, to the unevenness of the paving-stones in the courtyard of the Guermantes' mansion and the Baptistry of St. Mark's, it was because that being within me can only be nourished on the essence of things and finds in them alone its subsistence and its delight. It languishes in the observation by the senses of the present sterilised by the intelligence awaiting a future constructed by the will out of fragments of the past and the present from which it removes still more reality, keeping that only which serves the narrow human aim of utilitarian purposes. But let a sound, a scent already heard and breathed in the past be heard and breathed anew, simultaneously in the present and in the past, real without being actual, ideal without being abstract, then instantly the permanent and characteristic essence hidden in things is freed and our true being which has for long seemed dead but was not so in other ways awakes and revives, thanks to this celestial nourishment "While Deleuze's virtual ideas superficially resemble Plato's forms and Kant's ideas of pure reason, they are not originals or models, nor do they transcend possible experience; instead

they are the conditions of actual experience, the internal difference in itself. "The concept they [the conditions] form is identical to its object."[Desert Islands, p. 36.] A Deleuzean idea or concept of difference is therefore not a wraith-like abstraction of an experienced thing, it is a real system of differential relations that creates actual spaces, times, and sensations. [See "The Method of Dramatization" in Desert Islands, and "Actual and Virtual" in Dialogues. In fact it not as if the style and substance of science and philosophers have been entirely different. look at the conformality and consonance in the Gaussian Theory and Deleuze's statements on identity and the loss of it in the following passage…….That is, identity is a continuous process that finds its limit in exhaustion. As Deleuze will write, "…the logical relation of causality is inseparable from a physical process of signalling, without which it would not be translated into action. By 'signal' we mean a system with orders of disparate size, endowed with elements of dissymmetry; by 'sign' we mean what happens within such a system, what flashes across the intervals when a communication takes place between disparates. The sign is indeed an effect, but an effect with two aspects: in one of these it expresses, qua sign, the productive dissymmetry; in the other it tends to cancel it" (dr, 20). An actualized entity just is such a signal-sign system, remaining in perpetual communication with the virtual multiplicity it actualizes, constituting its own elements and identity. Such is Deleuze's account of what Badiou refers to as the operation of the "count-for-one"… an operation that is immanent to being itself and energetic in character. yet how are we to discern the local nature of ontological situations, the manner in which there is no global, overarching global situation in which all local situations are embedded as if parts of a whole? in part the above passage already responds to this question. if series must be brought into relation to communicate in order for causality to occur, and if the resonance of series is always a matter of chance, of a throw of the dice, like Lacan's objet a where no series enjoys primacy of model over copy, then all we have are local situations without a total situation. Unlike Heidegger's analysis of Dasein where Dasein is always being-in-the-world, Deleuze's simulacra or actualizations are always divergent, enjoying only local relations. Heidegger argues that in order to

understand the being of Dasein we must understand the manner in which Dasein is a being-in-the-world. By contrast, Deleuze's concept of multiplicity provides us with the principle of a new structuralism, of a new local ontology, that allows us to understand the immanent organization of a multiplicity without referring it to an embedding global space. as DeLanda so beautifully puts it in relation to the mathematician gauss, "…when gauss began to tap into these differential resources, a curved two-dimension surface was studied using the old Cartesian method: the surface was embedded in three-dimensional space complete with its own fixed set of axes; then, using those axes, coordinates would be assigned to every point of the surface; finally, the geometric links between points determining the form of the surface would be expressed as algebraic relations between the numbers. but gauss realized that the calculus, focusing as it does on infinitesimal points on the surface itself (that is, operating entirely with local information), allowed the study of the surface without any reference to a global embedding space. basically, gauss developed a method to implant the coordinate axes on the surface itself (that is, a method to implant the coordinate axes on the surface itself (that is, a method of 'coordinatizing' the surface) and, once points had been so translated into numbers, to use differential (not algebraic) equations to characterize their relations. as the mathematician and historian Morris Kline observes, by getting rid of the global embedding space and dealing with the surface through its own local properties, 'gauss advanced the totally new concept that a surface is a space in itself'" (intensive science and virtual philosophy, 11-12).according to the old Cartesian method, we can only outline the properties of a space by relativizing it to a global space in terms of which it is then mapped. by contrast, gauss is able to explore a space in terms of its intrinsic metric and organization as a local space, without referencing it to a whole of which it is conceptualized as a part. **The Image of Thought** "Image of thought "permeates both popular and philosophical discourse. According to this image, thinking naturally gravitates towards truth. Thought is divided easily into categories of **truth and error**. The model for thought comes from the educational institution, in which a master sets a problem and the pupil produces a solution which is either true or false. This image of the subject

supposes that there are different faculties, each of which ideally grasps the particular domain of reality to which it is most suited. In philosophy, this conception results in discourses predicated on the argument that "**Everybody knows..."** the truth of some basic idea. Descartes, for example, appeals to the idea that everyone can at least think and therefore exists. Deleuze points out that philosophy of this type attempts to eliminate all objective presuppositions while maintaining subjective ones. Deleuze maintains, with Artaud, that real thinking is one of the most difficult challenges there is. Thinking requires a confrontation with stupidity, the state of being formlessly human without engaging any real problems. One discovers that the real path to truth is through the production of sense: the creation of a texture for thought that relates it to its object. Sense is the membrane that relates thought to its other. Accordingly, learning is not the memorization of facts but the coordination of thought with a reality. "**As a result, 'learning' always takes place in and through the unconscious, thereby establishing the bond of a profound complicity between nature and mind".** Deleuze's alternate image of thought is based on difference, which creates a dynamism that traverses individual faculties and conceptions. This thought is fundamentally energetic and signifying: if it produces propositions, these are wholly secondary to its development. Philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; even tum tan tum;it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the vorticity of being merges with the positive use of the disjunctive synthesis, and

this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoin der for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for every hum of voices or a single silence for all the silences; a single vacuum for all the vacuousness consciousness should not be said without conjuring if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. The attributable to all states of bodies and states of affairs and the expressible of every proposition So univocity of consciousness means the identity of thé noematic attribute and that which is expressed linguistically and sense-fully. Univovity.We mean that it does not allow consciousness to be subsisting in aquasi state and but expresses in all pervading reality. (Reference: LOS: Deleuze) Deleuze sometimes spoke of getting up behind an author and "creating a monster", as if his reading methodology somehow consisted of a distortion of the philosopher's or author's thought. Certainly there is a case to be made for this with regard to the process of deterritorialization, where something is wrested from a territory and deterritorialized upon a new territory, like the animal paw that is deterritorialized from the earth and reterritorialize on the branch. however, it seems to me that the more interesting aspect of Deleuze's approach to other philosophers and art is not so much his "monstrous creations" (i seldom find them particularly monstrous), but rather their Gaussian or Riemannian style, where he explores them in terms of their own internal organization and metric, without reducing them to something alien such as history, society, biology (reductivism), or the signifier. What we find in Deleuze's approach to phenomena is a Gaussian technique. for instance, take Deleuze's books on cinema, his book on Francis bacon, or his study of sacher-masoch. In the first instance, Deleuze's carefully separates cinema from narrative and the signifier, studying it in terms of its specific organization pertaining to the production of images. in the case of bacon, Deleuze doesn't look for an underlying narrative or "meaning", but instead studies the manner in which bacon composes and

organizes his images and lines, both in terms of their production and actuality, so as to liberate a logic of sensation. finally, in approaching sacher-masoch, Deleuze allows Sacher-Masoch's novels to speak for themselves in terms of their desire and relation to pain, stalwartly refusing to reduce masochism to the complement of sadism. in each case, we have a local exploration of a "space" of multiplicities that is extremely precise. Sean Bowden the priority of events: **Deleuze's logic of sense sean Bowden, the priority of events: Deleuze's logic of sense, Edinburgh university press, 2011, 296pp., $40.00 (hbk), isbn 9780748643646 Reviewed by David Scott, Coppin state university and Gaussian spaces and multiplicities a note on Gaussian spaces and multiplicities posted by larval subjects under uncategorized.** What you "see" is not always what you "get." Many people mistakenly take their own visions literally without "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in the physical sense but in a psychological sense. The archetypes of the collective unconscious are arrayed behind the scenes of current worldwide conditions, of crisis and confusion. They mirror our own states back at us, whether we perceive them as such or interpret them plausibly or not. The noise of ordinary consciousness and beliefs drowns out the signal. Unconsciousness is the background of our ordinary awareness. Our organism is very much at the center of such effects. The organismic source is our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is (=) a meme-Scape. Pre-conceived concepts vie with (e&eb) structures, concepts with images. Like scientists who ignore (e) assumed truths, we leapfrog over (e&eb) our beliefs and personality deficits, claiming (eb) idiosyncratic imagination is (=) literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves (e) us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our

experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images. It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is (=) both. We inherit (e&eb) the structures that make our experience (=) what it is. But the structure itself is "empty," and each human culture "fills" it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about (e&eb) our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from

criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns,

nomenclature, and internal coherence of meaning and representation. n a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar some statements in Vedanta, it is not superfluous or redundant. It is evolutionary. Authetication determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication of the evolutionary process towards understanding is ever present in all, its thematic and discursive form. it is the relentless holding on to a theory this is despicable albeit the theory is not. Look at this paragraph. It is the 'I am the body/mind' belief that gives rise to the 'I am not (e) the world' belief. These two beliefs are (=) co-created."Consciousness projects (e&eb) the appearance of the mind, body and world by taking the shape of thinking, sensing and perceiving." "Attention is Consciousness with an object. When the object vanishes, attention simply remains what it always is, Consciousness.""There is (=) no purpose to (e) meditation. The purpose is already accomplished." "Everything that is experienced is experienced by, through, in and as Consciousness." "The seen cannot be separated from seeing and seeing cannot be separated from (e) Consciousness.""The Reality of any experience is not hidden in the appearance; it is expressed by the appearance." "Once we see that everything is Consciousness... Maya still dances, but it is a dance of love not seduction." There are some truly excellent sections

such as a long one taking, as a starting point, the following quotation from the painter Cézanne: "Everything vanishes, falls apart, doesn't it? Nature is always the same but nothing in her that appears to us lasts. Our art must render the thrill of her performance, along with her elements, the appearance of all her changes. It must give us a taste of her Eternity." Rupert begins: "That statement must be one of the clearest and most profound expressions of the nature and purpose of art in our era." And he goes on for some 14 pages to elaborate on this claim, examining the nature of the 'elements' Existence, appearance and Consciousness and their relationship

Some of the following concepts are modeled; some are not. For those which are not modelled kindly see one of the papers in the series. This note has the axiomatic predication to give a holistic view of the study. Here we talk of the characterstics of systems which satisfy the condition of cosmological constant or any other for that matter or variable or axiomatic predication. . There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit,

facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified

from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. Towards the end of circumvention of typing hundreds of superscripts and subscripts which would be a sardine squash and the concomitant operational difficulties, model is presented in piece meal of nine modules. Logarithms are to be taken in respect of those which incorporate more than one variable in bra-ket. Values of log (ab) and log (a+b) are readily available. Anti logarithm shall be taken at the end to predict, project, prognosticate the value of the variable. This is true for tensors, vectors and other variables too. Affirmational assertion, and explanation justification, statement of vindication, annotational commentary, and explicational glossary for each and every system changes and physical interpretation of results is one thing that is to be with earnest endeavour and feverish and febrile expediency. Editions never ever mean the same and identicality of parameters and this has been explicitly and unmistakably stated in the model apriori itself. Akin and analogous, cognate and concurrent, correspondent and congruent, comparable and complementary, synonymous and duple, tantamount and agnate, commensurate and correlative representation is only to highlight the importance and subterfuge the replication of work. It is my fervent solicitation to kindly bear with me for any lapses, notwithstanding the orchestrated efforts for a paper without any minor errors. Postulation predication, conclusive presumption, differential presuppositions, underscored decidedly axiomatic statement of the statement, equation form the bye word or the watch word in the aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement-argument, articulation, assertion, asseveration, choice of words, commentary and communication, declaration and definition, delivery-diction, elucidation-emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation tone and tenor of stratification. Now, how do find the reaction of systems to these singularities. You do the same thing a boss does for you. "Problematize" the events and see how you behave. I will resort to "pressure tactics". "intimidation of deriding report", or "cut in the increment" to make you undergo trials, travails and tribulations. I am happy to see if you improve your work; but may or may not be sad if you succumb to it and hang yourself! We do the same thing with systems. systems show conducive response, felicitous reciprocation or behave erratically with inner roil, eponymous radicality without and with blitzy conviction say like a solipsist nature of bellicose and blustering particles, or for that matter coruscation, trepidiational motion in fluid flows, or seemingly perfidious incendiaries in gormandizing fellow elementary particles, abnormal ebullitions, surcharges calumniations and unwarranted(you think so but the system idoes not!) unrighteous fulminations. Amount of emotion or affect that is appropriate or expected in an given form of interaction. Particular individuals and diffuse obligations (see c and d) are associated with affectivity, whereas contacts with many individuals (universalistic) in a bureaucracy may be devoid of emotion and characterized by affective neutrality. Affective neutrality may refer to self discipline and the deferment of gratification (eg Weber's spirit of capitalism). In contrast, affectivity may be associated with expressing emotions. Adams and Sydie also refer to affective neutrality being associated with ego control (p. 15). (Parsons: Wikipedia). So the point that is made here is "like we problematize the "events" to understand the human behaviour we have to "problematize" the events of systems to understand their behaviour. This statement is made in connection to the fact that there shall be creation or destruction of particles or complete obliteration of the system (blackhole evaporation) or obfuscation of results. Some systems are like "inside traders" they will not put signature atoll! How do you find they did it! Anyway, there are possibilities of a CIA finding out as they recently did! So we can do the same thing witn systems to. This is accentuation, corroboration, fortification, .fomentatory note to explain the various coefficients we have used in the model as also the dissipations called door. In the bank example we have clarified that various systems are individually conservative, and their conservativeness extends holisticallytoo.that one law is universal does not mean

there is complete adjudication of nonexistence of totality or global or holistic figure. Total always exists and "individual" systems always exists, if we donot bring Kant in to picture! For the time being let us not! Equations would become more ensorcelled and frenzied..... philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". In consciousness research, two rival sets of theories can be recognized: (A) Scientific material interpretations of consciousness are based on axioms that view consciousness in the context of highly advanced intentional processing of information in which subject-object relations evolve, and (B) humanistic interpretations of consciousness are based on axioms that view consciousness in the context of, say, "centered pulsations" that enable a conscious agent to act from his or her center of awareness. In this paper I will argue for the selection of axioms that favor humanistic interpretations of consciousness. The Pursuit of Autonomy Interdisciplinary Observations to Human Consciousness: Wautischer, Helmut (1993) the Pursuit of Autonomy. Interdisciplinary Observations to Human Consciousness: Social Neuroscience Bulletin, 6 (4). pp. 52-56.Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality; **CONSCIOUSNESS AND ITS UNIQUENESS June 12, 2012 at 8:40pm (See Deleuze Logic of sense, Wikipedia and Stanford encyclopedia for more details)** Atrocious contrivance and device, gratificational primogeniture, calamitous dodge and expedient, depraved, destructive, disastrous, execrable gambit and gimmick, great solace and succor, iniquitous, injurious, loathsome, low, maleficent, malevolent machination and maneuver, promethaleon of candor, frankness, honesty, honor, ingenuousness, innocence, openness, reality, sincerity, truthfulness, spiteful, stinking, ugly, unpleasant, unpropitious play and ploy, progenitor of jurisprudence and circumspection, racket and ruse, savvy, scam, stratagem, subterfuge, tactic and wile, forthrightness, honesty, truthfulness embodied and personified is the structural predisposition and dispensation of the complex. Nobody need have to take the Brahman-Anti Brahman agency to be a recalcitrant, repugnant, and refractory proposition. God gives and also takes away. That is the point made. ‘

And of course lesser gods, and God doth follow them and they follow God. That is why they are lesser gods. And wield the power on lesser mortals like us. When we mean gratification and deprivation we mean celestial Brahman Anti Brahman and Terrestial Brahman and Anti Brahman agencies producing such happiness or sadness by any means such as (blending- amalgamation, adulterant-adulteration, amalgam-amalgamation, blend- combination,

composite- compound, debasement-denaturant, fusion-hybrid, intermixture-reduction, amalgam- mixture, composite- compound, fusion-mishmash, aggregate- alloy, amalgamation-blend, commixture-composite, composition-compost, conglomerate-fusion, goulash-medley, mishmash-stew, blend-coalition, commixture-compound, federation-heating, integration-junction, melting-merger, synthesis-unification, assortment- combination, adulteration-alloy, assimilation-association, batter-blend, brewed combine, composite compound, concoction confection, conglomeration cross of money, orientation(dis), penury creation, murder, mayhem plunder, pillage, apocalypse, Armageddon in various forms and permutation and combination of Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness: with each action accounted and showing balance in the nature's general ledger as also actionlessness:e). All these are recorded in Nature's general ledger. Neuron DNA encapsulates all the actions passions, actions, interactions, transactions or the lack of it at various levels of individual, collective and cosmic levels. In celestial Brahman Anti Brahman ledger there is cessation of dualities. One more point to be made is if 5=3+2, then we can say 5 gormandizes 3 or 5 gobbles up 2. Quintessentially such a statement arises out of the difference of LHS and RHS term wise. It must be remembered despite all theoretical abstractions and generalizations be it proton mass or Higgs Boson, or topology of a space, it is notwithstanding the fundamentality of nature is just a number. Just everything in Physics and mathematics is just a number. Let us assume I buy rice in Bangalore (District) worth 5k.g. I can always divide 5k.g.s of rice in categories of Bangalore

(South), Bangalore (west), Bangalore (East) and Bangalore (South). On similar basis, there are billions of neutrons, protons and concomitant interactions which would satisfy the given numerical value. Example of gravity above makes things clear. Everything atleast is a function of time and someplace or some other coordinates, subordinates or superordinates. Mass of proton is 2.47 does not mean total mass of all protons is 2.47. It is this argument that is used in the classification scheme. By the very nature, self the unmonitoring but the witness consciousness agency identifies itself with something, be it a profession, an actor on the screen, when it does not have a identity. Such examples are legion, a Charlie Chaplin in "The Kid" or Raj kapoor in "Jagte Raho". These are people who just wanted to live and living without being put under surveillance every moment is the greatest part of their life. This is "Absolute Subjectivity" of Shiva. Such a state is always craving for identification for in pure consciousness all identities are lost. It is this "identitylessness" or having very minimal identity is what pervades most in India and this is due to expansion of individual consciousness. When that happens, be it education, or wealth, the "identity' becomes so well entrenched, they become the exact opposite of "pure consciousness"," impure consciousness" (Shakti). This is a state of "Absolute or relativistic objectivity". There was this case of a lady who went to such an extent of planning, she was sure that whatever she does must work out. When it did not, she lapsed in to "identity crisis". Such thing as this might happen in well developed countries where the competition is extremal and to be achieved at any cost. Realisation that one is not the doer but the witness consciousness is just like watching a movie in which one has acted. It is in this sense "advaita", "non duality' seems to be of cardinal and preeminent importance for the analytical treat ment of "identity' and "identity lessness", both happen in space time. "Multiverses" in this sense is also like a different film of an actor released at different theaters and it is in this sense we state that that every object in space time is a multiverse. You must have totally immersed in doing a "string theory" problem or watching a "film", oblivious of surroundings. Essentially creation (space time) also is what is happening in mind and honed in to you billion times. Identification of anything with a designation or signification is one that leads to

association within spacetime, which would only have pathological consequences and detrimental ramification as stated above or positive implications of "happiness" because of the identity itself. Search or 'show" has to go on! It gives a sense of consummation. a watching pot never boils; donot go to check nature; you are disturbing it; it will never show its true colours bhakti and Virakthi, gratification and deprivation, raga and dwesha is essential to live in space. It is the quantum interest thereof that produces energy to live .gratification and deprivation which form the bastion, pillar, post and stylobate and sentinel of all human actions depends on tams, rajas, and sattva of individual or insentient object. Verily tamas, rajas and satva are the characterstics of space itself, nay the behavioural pattern and attitudinal orientation. There are peaceful places like Vienna or Sringeri and there are violent places like others. Space is warped and woofed by the Akshara, the quantum information of primordial syllable Om. cosmic general ledger is the one vibration of which is brahman; quantum information possessions thereof form the sentinel of all transactions (Gargi to Yajnavalkya: my interpretation)adhocism, nonseqientiality of events, anti sequentiality of events, simultaneity of phenomenon, concurrent happenings , improbability of phenomenon happening itself are characterstics of brahman-anti brahman (Terrestial and celestial) are the cause of non predictability of happenings like in black hole, or in the areas of operation of brahman-anti brahman (Terrestial and celestial). brahman AntiBrahman actions, interactions, decisions, circumspection, adhocism, satisfies sensitive to initial conditions; topologically mixing; dense periodic orbits, the axiomatic predications and postulation alcovishness of chaos theory. Like gravity is a characteristic of space, so is tamas, rajas, sattva, with spaces knotted together by quantum information (for models see below) .accidents are preplanned actions of Terrestial and celestial Brahman AntiBrahman. All actions are perpetrated by them, be it train on fire or death. Belief systems play an important role in the determination of tamas, rajas, and satva state of human beings. There is neither cause nor effect. All actions are performed by Terrestial and celestial Brahman anti Brahman and it is off the record. psychic energy is negative equivalent to –mc^2, so that second law of thermodynamics is satisfied on a holistic basis vis-à-vis material energy

ParaBrahaman is infinite individual, collective and cosmic consciousnesses are infinite (concomitant general ledgers) for they are in the evolution mode with respect to each other in that order. Later emanates from the former as transactions emanate from the conservative state of assets and liabilities. Is not zero true for infinite conservative systems inside space and time? zero is the only consummation, consolidation and conservation. if fifty out of seventy members come out of the room at whatever time you come out, talk about the grand limbo of oblivion and hibernation of lifeless world you are going to and if you are not a celebrity, then it is not coincidence, it is coordination. Nature is playing the same trick on you. That a generalised law is there and you can measure does not make things less good albeit recalcitrant, refractory and rogue particles occur time and again, by and large. Entering in to all living beings through his constant part on the new moon night, he is born therefrom next morning. Absolute ParaBrahaman (zero) projects himself again and again on the screen of individual consciousness which appears of depairing despondent world. Some where the balancing act is done on karma. cosmic general ledger reflects upon itself to keep all transactional ties conservative (Brihadarnkayaka Upanishad: 1.5 14) for models see next paper through its divine power ( i interpret man creates god; it is attribution and ascription of various Shaktis would the formation of god takes place like space and time) the self assumes (eb) the form of a deity man can contemplate and venerate, even though siva, the pure subject, can never in fact be an object of meditation (adhyeya). Until we realise our true identity with Sankara, he is worshipped and conceived to be a reality alien to ourselves. While we are in the realm of creation, he too is a creation or mode or appearing of the absolute, manifest to us in meditation, through his freedom as an eternal, omniscient being. There is no gulf between the created and the uncreated creator: nothing in reality, although an object of knowledge, ceases to be Siva: this is the reason why meditation [on this or that aspect] of reality bestows its fruit. The world of the senses and mind appears to the well awakened (suprabuddha) as a theophany, an eternal revealing of god in his creation. The doctrine of vibration declares that "there is no state in word, meaning or thought, either at the beginning, middle or end, which is not Siva." 59 to utter any word is, in

reality, to intone a sacred formula.every act is a part of Siva's eternal cosmic liturgy, every movement of the body a ritual gesture (mudras), and every thought, god's thought. By what path are you not attainable? What words do not speak of you? in which meditation are you not an object of contemplation? What indeed are you not, o lord? 61 spandashakti, which accounts for the appearing of all things, is also the means by which deity in its many varied forms appears to man. Ksemaraja concludes: the ultimate object of worship of any theistic school differs not from the Spanda principle. The diversity of meditation is due solely to the absolute freedom of Spanda. 62 Sankara is not only the supreme object of devotion; as the static polarity of the absolute, he is the inner reality which holds together its Siva and Sakti outer manifestations. Phenomena are patterns of cognitions projected onto the surface of self-luminous siva-consciousness. There they become apparent, directly revealed to consciousness according to their manifest form. Siva is accordingly symbolized as the ground or surface of awareness, smooth and even like a screen (samabhittitalopama). 64 inscribed on this screen (kutfya) are the countless manifest forms which appear within it rendering it as diverse and beautiful as a fossil ammonite (falagrama). 65 Siva is the sacred ground upon which the cosmic mandalas are drawn, the absolute surface of inscription which bears the mark (cihna) of the universe. Abhinava writes: the variety of this world can only be manifest if the highest lord, who is essentially the pure light of consciousness, exists; just as a surface is necessary for a picture. if external objects were perceived in isolation then, because 'blue' and 'yellow', etc., are self-confined and the perceptions [we have of them] refer to their objects alone and so are insentient, mute and dumb in relation to one another . . . how would it be possible to be aware that an object is variegated? but just as depths and elevations can be represented by lines on a smooth wall, and we perceive [a female figure and think], similarly it is possible to be aware of differences in the variegated (contents of experience) only if all the diverse perceptions are connected together on the one wall of the universal light of consciousness. 66 Siva is the perfect artist who, without need of canvas or brush, paints the world pictures. The instant he imagines it, it appears spontaneously, perfect in every respect. The colours he uses are the varying shades

and gradations of his own Spanda energy and the medium his own consciousness. Intractable disaffect, Mutinous obdurate, Rebellious recalcitrant, Seditious refractory, Insubordinate insurgent, antagonistically antipodean, discrepant incompatible, Contumacious coadjutant, Conflicting converse, Unyielding obdurate, the antagonistic and anachronistic dispositional ties of Anti Brahman are in fact the sine qua non of Brahman only (See Vishwaroopa of Mahavishnu, which contains the polarised positives on RHS and negatives on LHS). The universe is coloured with the dye of its own nature (svabhava) by the power of Siva's consciousness (a//). Rajanaka rama says: homage to him who paints the picture of the three worlds, thereby displaying in full evidence his amazing genius {pratibha)\ to Sambhu who is beautiful with the hundreds of appearances laid out by the brush of his own unique, subtle and pure energy. 68 analogously, at the microcosmic level, all the cognitions and emotions, etc., which make up the individual personality form the outward flow of essentially introverted consciousness. They are specific pulsations (vise$aspanda) or aspects of the universal pulsation (samanyaspanda) of pure t consciousness. at the lower level, within the domain of maya, they represent the play in the fettered soul of the three primary qualities (gunas) or feeling-tones' which permeate to varying degrees his daily experience. These are: 1) sattva — the quality of goodness and luminosity which accompanies blissful experience both aesthetic and spiritual. 2) Rajas — the passion or agitation which oscillates between the extremes of 'light' and 'darkness' and characterises inherently painful experiences. 3) Tamas — the torpor and delusion which accompany states of inertia and ignorance. 69 the liberated soul recognises that these three are the natural and uncreated powers of pure consciousness. for him they are manifest respectively as: 1) Sankara 's power of knowledge (jnana) — the light of consciousness (prakasa); 2) the power of action (kriya) — the reflective awareness of consciousness (vimarsa)\ 3) the power of maya — which does not mean here the world of diversity, but the initial subtle distinction which appears between subject and object in pure consciousness. 70(doctrine of vibration: mark S.G. Dyczkowski) .to "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new

concepts and meanings that could not be included in the old system. at the heart of western metaphysics, for example, derrida finds the opposition between "speech" and "writing." this binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior in derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. The postmodern turn in philosophy: theoretical provocations and normative deficits by Steven best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html. You cannot see beyond consciousness for that is the limit of your storage of information or the basis for understanding further information. Turiya is empty bucket; it is from which prajna the measure of  arises; turiya measures the number of quantum information stores in other three states of waking, dreams, and dreamless deep sleep (Manduka Upanishad; my interpretation for models next paper).let there be path where angels tread: it is a force never born never dies the brahman an anti brahman verily there are parallel universes in (eb) in the ambit of space and time; the society of sinners and the sinned with the grammar of evil; parallel universes contain(e) multiverses in them; all physics, philosophy,  break down(e) in these universes; astronomical gratification ,deprivation, glorification, mortification, projections, axes, dimensions, rotations, foldings, concurrences, motions, shocks, are the axiomatic predications of multiverses and parallel universes; death is when self (the witness consciousness; register of individual general ledger) leaves the body; ostensibly parallel seem to have (e) an objective; parallel universes donot (e) have objectives; their paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common

attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, affection- appreciation, awareness- discernment, emotion-feeling, gut reaction, heart insight, intuition, judgment, keenness, perceptiveness, rationale, sensation, sense, sensitiveness, sensitivity, sentiment, susceptibility, taste, vibes, admiration, aesthetic sense, affection, appraisal, assessment, attraction, awareness, cognizance, commendation, comprehension, enjoyment, esteem, estimation, grasp, high regard, knowledge, liking, love, perceptional realization, recognitional regard, relish, respect, responsiveness, sensibility, sensitiveness, sensitivity, sympathy, understanding, valuation, acknowledgment, airing, baring, , confessional defenselessness, denudation- denunciation, disclosure, display, divulgence-exhibition, hazard, introduction, jeopardy, laying open, liability, manifestation, nakedness, openness, peril, presentation, publicity, revelation, risk, showing, susceptibility , susceptiveness,susceptivity, unfolding, unmasking, unveiling, vulnerability, vulnerableness, action, affection, appreciation and ardor, behavior, capacity, compassion, concern, cultivation and culture, delicacy, discernment and discrimination, emotion, empathy, faculty, fervor, fondness, heat, imagination, impression and intelligence, intensity and intuition, judgment, keenness- palpability, passion, pathos, pity, reaction, refinement, sensibility, sensitivity, sentiment, sentimentality , sharpness, spirit, sympathy, tangibility, taste, tenderness, understanding, warmth, fond memories, hearts and flowers, homesickness, longing, pining, reminiscence, ,schmaltz, sentimentality wistfulness, yearning satiation is only ostensible; their job is to get the job done; their business is to mind other's business ;all laws of physics and finance breakdown; it is antilaw,piscatorial Piratish and bubonic bucaneerishness; there is free transactions between in multiverses; when reflected on the screen of consciousness they appear as extant existential universe. Transactions are astronomical and gratification

deprivation, glorification and mortification incalculable; they have no barriers of region, religion, language which they create for others. theirs objectives are: corruption criminality, debasement debauchery, degradation evil immoral lewdness licentiousness, perversion sinfulness, vice viciousness, vileness wickedness, improbity corruption, atrocity decadence, degeneration degradation, depravity evil, immorality impurity, infamy iniquity, looseness lubricity, perversion profligacy, sinfulness turpitude,vice viciousness, vulgarity wickedness,disgrace bad reputation, abasement abuse, baseness black eye, blemish blur, brand comedown,contempt contumely, corruption culpability, debasement debasing, defamation degradation, derision disbarment, discredit disesteem, disfavor dishonor, disrepute disrespect, humbling humiliationignominy ill repute, infamy ingloriousness, meanness obloquy, odium opprobrium, pollution ,put-down reproach, scandal scorn, slander slight, slur spot, stain stigma, taint tarnish, turpitude venality, enormity Horribleness , abomination atrociousness, atrocity crime and the positive aspectionalities thereof .let there not be any incursions where evil mercenaries in the garb of mendicants thrive with adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatious gainsaying or conformingly pugnacious Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, execrable and goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness . statement for

Mephistophelean mercenaries philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality. Negative energy stabilising wormholes are considered candidates for (er) darkmatter. An analysis of negative energy might be a useful ingredient of the thoughts of travel at speeds more than the velocity of light. Peirce explains that an "imperfect continuum" possesses "topical singularities, or places of lower dimensionality" which break continuity, while "the parts of a perfect continuum have the same dimensionality

as the whole." (CP 4.642) Circular continua without singularities are time and space. (CP 6.210-212) By using math as his lens, Peirce envisioned singularities or lower dimensional holes in the imperfect cosmologic continuum, while maintaining that a probabilistic law of chance exists that entropy decreases **(Wikipedia).** On the side of philosophy**,** quite explicitly, Deleuze provides a critique of Husserl's "sleight-of-hand" change from a static categorical phenomenology to one supposedly more dynamically constitutive of reality. This a clear cut **chagrined circumvention**, **contraventional curbing**, **defeated disgruntlement**, **dissatisfactional downer**, **dragged failure**, **fizzled foil**, **grievance hindrance**, **impedimental irritation**, letdown, nonfulfillment, nonsuccess, obstruction, old one-two, **resent mental setback**, unfulfillment of the earlier theories makes Sartre contend that human existence is a conundrum whereby each of us exists, for as long as we live, within an overall condition of nothingness (nothing-ness)—that ultimately allows for free consciousness. But simultaneously, within our being (in the physical world), we are constrained to make continuous, conscious choices. It is this dichotomy that causes anguish, because choice (subjectivity) represents a limit on freedom within an otherwise unbridled range of thoughts. Subsequently, humans seek to flee our anguish through action-oriented constructs such as escapes, visualizations, or visions (such as dreams) designed to lead us toward some meaningful end, such as necessity, destiny, determinism (God), etc. Thus, in living our lives, we often become unconscious actors—Bourgeois, Feminist, Worker, Party Member, Frenchman, Canadian or American—each doing as we must to fulfill our chosen characters' destinies. However, Sartre contends our conscious choices (leading to often unconscious actions) run counter to our intellectual freedom. Yet we are bound to the conditioned and physical world—in which some form of action is always required. This leads to failed dreams of completion, as Sartre described them, because inevitably we are unable to bridge the void between the purity and spontaneity of thought and all-too constraining action; between the being and the nothingness that inherently coincide in our self. Deleuze insists on the absolute non-resemblance between what conditions (metastable, pre-individual plane of singularities) and what is conditioned

(actualized or the individuated 'thing'). Only then do "the conditions of the true genesis become apparent," Deleuze writes (Logic of Sense 105). Without addressing this issue, Deleuze's borrowings from Simondon, Lautman and structuralism remain untethered. Because of Bowden's decision to completely sidestep Deleuze's overt negotiations with phenomenology, his otherwise sophisticated commentary suggests a pedagogical narrowness. If, the omission of phenomenology strengthens Bowden's focus on the ontological priority of the event over substantial ontologies via the genealogy of figures and movements Deleuze outlines, it also incompletely articulates the problem of the sense-event as Deleuze formulates it. We have lost any idea of whom Deleuze takes to be his primary adversary, against whom he stakes out his position. This worry is most clearly manifested in Bowden's description of Deleuze's ontological prioritizing of events as offering a "transcendental ontology." For Bowden Deleuze's transcendental ontology defines events that "are ontologically prior to substances, 'all the way down'" (82). But Bowden provides a second definition: it is that the priority of events over worldly individuals means that the world of events "is not something external to the conditions of knowledge" (69). **SEAN BOWDEN the Priority of Events: Deleuze's Logic of Sense Sean Bowden, the Priority of Events: Deleuze's Logic of Sense, Edinburgh University Press, 2011, 296pp., $40.00 (hbk), ISBN 9780748643646. Reviewed by David Scott, Coppin State University Weltschmerz, agony, apprehension, blues, depression, dread, mid-life crisis, misgiving, nervousness, uneasiness, angst were expressed with such effulgent words…… aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement, argument, articulational assertion, asseverational choice of words, commentary and communication, declaration and definition, delivery, diction, elucidation, emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation about the human predicament is brought out in Being and Nothingness. Being and Nothingness offers a critique of Sigmund Freud's theory of the unconscious, based on the claim that consciousness is essentially self-conscious. Sartre also argues that**

**Freud's theory of repression is internally flawed. According to Sartre, in his clinical work, Freud encountered patients who seemed to embody a particular kind of paradox—they appeared to both know and not know the same thing. In response, Freud postulated the existence of the unconscious, which contains the "truth" of the traumas underlying the patients' behavior. This "truth" is actively repressed, which is made evident by the patients' resistance to its revelation during analysis. Yet what does the resisting if the patients are unaware of what they are repressing? Sartre finds the answer in what Freud calls the "censor". "The only level on which we can locate the refusal of the subject," Sartre writes, "is that of the censor.** The Holographic Principle, as conceived in current physics, applies to fields, and perhaps, even to more elementary entities, called strings and branes. These more elementary entities remain hypothetical at the time of this writing. There are many layers between the level of fields and that of neurons. We know that neurons are surrounded by a surface phospholipid membrane, which supports an electrochemical process that is fundamentally necessary to human experience. As the information supported by the membrane surface according to the Holographic Principle is proportional to its area, we should expect to find that the neural correlates of consciousness most clearly associated with those parts of the neuron that have the have a high proportion of neuronal surface area, and, in particular, receptive surface area. Dendrites account for an average of about 90% of the neuron is receptive surface area (Wong, 2002), and so would be expected to be the most important neural structures in information processes. We find this to be true **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science (Wikipedia).** Discordant-discrepant, dissonant-hard-line, implacable, incompatible, incongruous- inconsistent, inexorable-inflexible, inharmonious-intransigent, opposed, reluctant, unappeasable, uncompromising, unfriendly some of these expatiations and enumerations might sound, a more generalised interpretations is given by a Vedantin. Look at this arresting passage: The teacher continued to be silent. When addressed a second and third time he said: "I am teaching, but you do not follow. The Self is silence." The undetermined and unthinkable character of the Brahman is

a consequence of the absolute's eternal and immutable nature. To concede the existence of a real universe is, from the Vedantin's point of view, to posit the existence of a reality apart from the Brahman. Nor can we simply identify a real universe with the absolute unless we are prepared to compromise its unchanging, absolute status. The criterion of authenticity is immutability. Reality never changes; only that which is less than real can appear to do so. Reality is constant in the midst of change. What this means essentially is that there is change although nothing changes. This impossible situation is reflected in the ultimate impossibility of change itself. That which does not exist prior to its changing and at the end, after it has changed, must be equally non-existent between these two moments. Although the world of change appears to be real, it cannot be so. Change, according to the Vedantin, presupposes a loss of identity. Reality cannot suffer transformation; if it were to do so, it would become something else and the real would be deprived of its reality. The immortal can never become mortal, nor can the mortal become immortal. The ultimate nature of anything cannot change. Change of any sort is merely apparent (vivarta); the world of change and becoming is a false super- imposition (adhyaropa, adhyasa) on the absolute. In cosmic terms, the mistake (bhranti) consists of the supposition that the real Brahman is the unreal universe and the unreal universe is the real Brahman. In microcosmic terms, it is the mistake of falsely conceiving the body, mind or even one's personality to be the Self. In the same way as the image of a snake is falsely superimposed on a rope, similarly the universe is falsely projected onto the real substratum, the Brahman. Ignorance is not merely a personal lack of knowledge, but a cosmic principle. As such it is called "Maya," the indefinable factor (anirvacaniya) that brings this mistake in identity about. The reality status of this cosmic illusion is also indefinable: on the one hand it is not Brahman, the sole reality; on the other hand it is not absolutely non-existent like a hare's horn or the son of a barren woman. Brahman is the source of world appearances only in the sense of being their unconditioned ground or essential nature. The universe is false not because it has no nature of its own but because it does have one. Just as the illusion of a snake disappears when one sees that it is nothing but a rope, similarly cancellation (badha) of the empirically real occurs

when the absolute reality of the Brahman is realised. Thus, according to Vedanta, appearance implies the real, while the real need not imply appearance. To appear is essentially to appear in place of the real, but to be real is not necessarily to appear. All things -exist because the absolute exists. It is their Being. Thus the very existence of phenomena implies their non-existence as independent realities.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^\infty \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9 \left( U_9 - V_9 \int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9 \right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9 \int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$then given the constraints of $f(T_9), (A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 -$$

$$V_9 M_9 -]e^{-T_9\lambda_9} = 0$$

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9 Z_9 \sin(Z_9 T_9) - (A_9 M_9 + Y_9) + (A_9 M_9 + X_9)\cos(Z_9 T_9) = 0$$

$$(A_9 M_9) Z_9 - M_9 Z_9 \cos(Z_9 T_9) + (A_9 M_9 - X_9)\sin(Z_9 T_9) = 0$$

Quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High

Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) **Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP** This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) **Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP** In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference: Journal of Physics: Conference Series 70 **Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows:** System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol

Where $A_9 = U_9 - V_9$Using (A2.3), substitute for $\boldsymbol{Z}_9^2$ and $Z_9$in (A2.2):

$W_9^2 (X_9^2 - 3 M_9 A_9 X_9 - M_9^2 (X_9 - Y_9) + 2 M_9^3 A_9 + 2 A_9^2 M_9^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_9$=1-cos ($Z_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (X_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ X_9 - Y_9/2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \leq W_9 \leq 2$ *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (X_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\,M_9)^2\,(2A_9 M_9 - (X_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9 M_9^2 - 3M_9 X_9\right) + W_9\left(3M_9 A_9^2 - 2A_9 X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9+Y_9\right)$$

$$=2M_9^2W_9^2(M_9+\mathrm{A}_9)+M_9A_9^2W_9+\left(2A_9M_9-X_9+Y_9\right)(A_9W_9+M_9W_9^2$$

$$+(2-W_9)((2M_9X_9+M_9Y_9)W_9+A_9(X_9+Y_9)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range 0 $W_9$ 2.Defining

$Q_9\left(W_9,A_9\right)=a_9W_9^2+c_9$ *where* $a_9b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

To Heidegger writings we shall compose the model. Following is not still written model able form. For one of the models please see any one of the papers under the appellation: Ïn Excited Reveries

In his later writings Heidegger drops the term 'phenomenon' in preference for the verbal form 'phainesthai' in order to emphasize even more the actuality or presentational property of Being. Explaining this new form of the term he writes: "Being disclosed itself to the ancient Greeks as 'physis'. The etymological roots 'phy-' and 'pha-' designate the same thing: 'phyein', the rising-up or upsurge which resides within itself as 'phainesthai', lighting-up, self-showing, coming-out, appearing-forth." Heidegger contrasted his notion of phenomenon with semblance (Schein) and with appearing (Erscheinung). In the case of semblance a thing can show itself as that which it is not, as when fool's gold shows itself to be gold. The ancients always allied semblance with non-being. Heidegger points out, however, that semblances are grounded in

showings, and so does Abhinava. Both Heidegger and Abhinava consequently maintain that all semblances have a real basis and are to be treated as instances of phenomena along with the so-called real showing or manifestation of non-deceptive objects. So Heidegger states that: 'how- ever much seeming, just that much being'. Thus self-showing or appearing defines Being as phenomenon, but this definition of Being is as yet incomplete. Being is not only self-showing but 'logos' which Heidegger explains means 'discourse' (Rede) in the sense of 'apophansis': 'letting-be-seen'. Phenomenology, which according to Heidegger is the only correct study of Being, means 'letting-be-seen-that-which- shows-itself. This is true of Saiva Paramadvaita as well. The reality of the world demands recognition; we are forced to accept the direct presentation of the fact of our daily experience. As Abhinava says: "if practical life, which is useful to all persons at all times, places and conditions were not real, then there would be nothing left which could be said to be real" A thousand proofs could not make 'blue' other than the colour blue. The reality of whatever appears in consciousness cannot be denied. Objects appear; they do not cease to do so by a mere emphatic denial. The manifestation of an entity in its own specific form is a fact at one level of consciousness; it is real. The appearing of the same entity in the same form but recognised to be a direct representation of the absolute is also a fact, but at another level of consciousness. It is no more or less real than the first. 'As is the state of consciousness, so is the experience,' says Abhinava. Although the nature of the absolute is discovered at a higher level of consciousness, The Doctrine of Vibration nonetheless it presents itself to us directly in the specific form in which we perceive things; otherwise there would be no way in which we could penetrate from the level of appearing to that of its source and basis. Abhinava writes: Real is the entity iyastu) that appears in the moment of direct perception (sak$atkara), that is to say, within our experience of it. Once its own specific form has been clearly determined one should, with effort, induce it to penetrate into its pure conscious nature. All things are known to be just as they present themselves. The concrete actuality of

being known (pramiti), irrespective of content, is itself the vibrant (spanda) actuality of the absolute. Liberating knowledge is gained not by going beyond appearances but by attending closely to them. "The secret," Mahesvarananda says, "is that liberation while alive (jivanmukti) is the profound contemplation of Maya's nature." No ontological distinction can be drawn between the absolute and its manifestations because both are an appearing (dbhasa), the latter of diversity and the former of 'the true light of consciousness which is beyond Maya and is the category Siva'. Those who have attained the category of Pure Knowledge above Maya and have thus gone beyond the category of Maya, see the entire universe as the light of consciousness . . . Just as the markings [on a feather] are nothing apart from the feather, the feather [is nothing apart from] them, similarly, when the light of consciousness is manifest, the whole group of phenomena is manifest as the light of consciousness itself. Within the sphere of Maya, every entity's 'own nature' (svabhdva) corresponds to its specific manifest form. Accordingly it is defined as that which distinguishes it from all else and from which it never deviates. Above the sphere of Maya, that is, above the level of objectivity, is the domain of the subject. At this level, everything is realised to be part of the fullness of the experience r and hence no longer bound by the conditions which impinge on the object. Here the part is discovered to be the whole, that is, consciousness in toto. In this sphere beyond relative distinctions, the yogi realises that (all) the categories of existence are present in every single category. m The yogi experiences every individual particular as the sum total of everything else. He recognises that all things have one nature and that every particular is all things. This is the 'essence' (sard) or co-extensive unity (samarasya) of all things. Integral Monism of Kashmiri SaivismWe have established that reality is manifest according to how [and the degree in which] the freedom of consciousness reveals it and that [this freedom] is the womb of all forms. Just as 'sweetness' is present in its entirety in every atom of the sugarcane, so each and every atom [of the universe] bears within itself the emanation of all things. This is

the level of consciousness in which the absolute reflects on itself realising to its eternal delight and astonishment (camatkara) its own integral nature. The reality of the world of diversity is not denied, but experienced in a new mode of awareness free of time and space in the eternal omnipresence of the Here and Now. [Phenomenal forms of awareness] such as 'this [exists]', born of the colouring [imparted to the absolute] by the limitations engendered by the diversifying power of time {kalakalana) also emanate within the Supreme Principle. There [at that level], Fullness {purriata) is the one nature [of all things] and so everything is omnipresent; otherwise, associated with division (khancjlana), the Fullness [of the absolute] would not be full. The content of absolute consciousness consists of diverse appearings (abhdsa) which, because they are manifest through it in this way, do not compromise the wholeness of consciousness. Everything we perceive is a momentary collocation of a number of such manifestations which combine together like 4 a row of altar lamps' (dipavalT) to form the single radiant picture of the universe. The individual objects which constitute the universe are specific collocations of such 'atomic' appearings. Together they form a single unified particular which appears according to its own defining features (svalaksand). A jar, for example, consists of a number of appearances such as 'round', 4 fat\ 'earthen', 'red', etc., which together discharge a single function (arthakriya), in this case, that of carrying the appearance 'water'. They unite with each other much as the scattered rays of a lamp come together when focused, or as the various currents of the sea together give rise to waves. Atomic appearings can combine in any number of ways, provided that they are not contrary to one another as established by the dictates of natural law (niyati). An appearance of 'form', for example, cannot combine with that of 'air'. Insofar as they share a common basis (samanyadhikaranya), a given cluster of appearances appears as a single whole. This common basis is the most prominent member of the group; the appearance 'jar' is such in the example quoted above. Any one appearance in a cluster may assume a more important or subordinate role.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_9 \square - b_9 \square 0$$

And

$$b_9^2 \square 4a_9 c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9 in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\ X_9$ ,

$Let\ Y_9 = \ y_9\ X_9$And try the substitution $A_9 = X_9/(2 + A_9\ y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + \ n_9 y_9)$ we have

$$W_9^2\ y_9\left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^2\left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9\,X_9\left[\frac{X_9\left(1+n_9y_9\right)}{M_9^2\left(2+n_9y_9\right)^2}+2\,y_9\left(n_9-1\right)+\frac{X_9^2\left(9+y_9\right)}{M_9^2\left(2+n_9y_9\right)}\right]=0.\ (A2.12)$$

Since $n_9$ and $A_9$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9=\left(2+n_9y_9\right),$ *the condition*$(A2.99?)$*can be written as*

$$\left[\frac{X_9\left(E_9-9\right)}{E_9^2}\right]^2\left[\frac{X_9^2}{M_9^2}-4\,y_9\,E_9^2\frac{X_9}{M_9^2}+4y_9\left(n_9-2\right)E_9^3\right]\ \text{[illegible]}0.$$

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9/W_9^2)$ can be expressed as

$$2\,E_9\left[y_9E_9+\left(y_9\,E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}\right]>\frac{X_9}{M_9^2}>2E_9\left[y_9E_9-y_9E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}]$$

dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique

for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) **Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP** Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) **Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP** Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a

cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) **Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP** Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service

placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) **Dynamic Service Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP** Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by $H*(X,(\wedge*TX)\otimes g)$, where X is the target hyperk\"ahler manifold. We further give a mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) **A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van**

**Raamsdonk1 and Nathan Seiberg2** study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. **Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP** Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy **The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP** Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the

present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. **ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP** Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition. Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb) the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537

© 2016 American Physical Society **Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP** Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant ($\Delta\alpha/\alpha$) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and $\Delta\alpha/\alpha$ to (e&eb) this conclusion. However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts $z\lesssim 3$, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the $\Lambda$ cold dark matter model value ($j\Lambda CDM=1$); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in $\alpha(z)$, linked to (e&eb) a linear coupling ($\zeta$) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order $|\zeta|\sim 10-4$, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in $\Delta\alpha/\alpha$ (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society **The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP** Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation

rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void

Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems, authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society **Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP** Authors measure the star formation properties of (e) two large samples of

galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the **star formation** rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a **function of environment**. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]**.** Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size $\sqrt{n} \times \sqrt{n}$? The problem is that $\sqrt{n}$ time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time $O(\sqrt{n})$ for $d \geq 3$, or $O(\sqrt{n}\log 5/2n)$ for $d = 2$. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search

tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ∗This work was mostly done while the author was a **PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship and by CIAR**. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 **Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP** Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can

lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 **Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP** Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. **Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP** Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e)

sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) **Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP** The relationship between the foundations of mathematics and (e&eb) physics is a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb) flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling

field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) **Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP** The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar

and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph] (or arXiv: 1503.05600v1 [math-ph] for this version) **Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP** This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs

mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) **Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP** Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and

derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to **the fine structure constant**. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) **Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP** Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor goes to 0 as the (e&eb) time goes to 0, the

big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph] (or arXiv: 1306.4613v1 [quant-ph] for this version) **Effects of gauge theory based number scaling on geometry Paul Benioff STOP** In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x,

outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph]

(or arXiv: 1211.3381v2 [quant-ph] for this version) **Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP** Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value,

\int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances, the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) **Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP** The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include

(e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) **Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP** Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) **New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP** This work is based on (e) the field of reference frames based on (e) quantum

representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects:Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc) Journal reference:Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) **A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul B**

**ABSOLUTE STABILITY**

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 ? $W_9$ ? 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9-9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2-2y_9\frac{(E_9+n_9^2y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right)-4y_9(E_9-n_9)$$ (A2.14)

And has the solution

$$\frac{X_9}{M_9^2}>\frac{E_9}{(E_9-9)}[y_9(E_9+n_9^2y_9)+[y_9^2(E_9+n_9^2y_9)^2+4y_9(E_9-9)(E_9-n_9)]^{9/2}]$$

Given $E_9>n_9$this lower limit to $X_9/M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9/M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9<n_9$ $(ie.,2A_9M_9^2<X_9-Y_9)$then (A2.9)and (A2.8)show that $Q_9$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2}>2E_9[y_9E_9-(y_9E_9(9+y_9)(E_9-n_9)]^{9/2}]$$

And

$$E_9>n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2/2M_9(X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$ and $X_9$ and $Y_9$ by $W_9^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

**Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability

of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e) long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007 STOP. David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer P** examine the possibility of coherent reversible information transfer between solid-state superconducting qubits and (e&eb) ensembles of ultracold atoms. Strong coupling between these systems is mediated by (e&eb) a microwave transmission line resonator that interacts near resonantly with (e&eb) the atoms via (e&eb) their optically excited Rydberg states. Solid-state qubits can then be used to (e) implement rapid quantum logic gates, while (e) collective metastable states of (e) the atoms can be employed for (e) long-term storage and (e&eb) optical readout of quantum information. **Reversible state transfer between superconducting qubits and (e&eb) atomic ensembles David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer Phys. Rev. A 79, 040304(R) – Published 14 April 2009 STOP.** Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures have been proposed. An

attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. **Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier** demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in (eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. **Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP.** In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown

to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , **Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally** frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. **Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg** propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states

into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. **Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011 http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan** proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 **Qubit transfer between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller** propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. **Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys.**

**Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich** report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. **Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi** investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. **Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev** investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer.

Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. **L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP.** In recent years **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P have** investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P** apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys. Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. **Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S.**

**Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP.** Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert **Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe** realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.**220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2** [quant-ph] for this version) **Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP.** Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. **Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker** describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue

from over (e&eb) a thousand plant and animal samples. **NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP.** Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this behavior, enhancing and accelerating (eb, eb+) the growth of correlations. **Michael Kastner** gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. **Michael Kastner** tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) **Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP.** A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. **Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E. Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91,**

**075435 – Published 26 February 2015 STOP.** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still

problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to

(e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP** Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that

(eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035

**Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime**

**C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013**

**STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and

cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP** Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex

Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) **A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen**

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9 and$ separate (A3.1)into real and imaginary parts

$$Z_9^2 + \left(M_9 + V_9\right) Z_9 \sin\left(Z_9 T_9\right) - \left(U_9 M_9 + Y_9\right) + (M_9\left(U_9 + V_9\right) - X_9)$$

$$\mathrm{X}\cos\left(Z_9 T_9\right) - 2M_9 V_9 cos^2\left(Z_9 T_9\right) + M_9 V_9 = 0$$

$$(M_9 + U_9)\, Z_9 - (M_9 + V_9)\, Z_9\ cos\left(Z_9 T_9\right) + \left(M_9\left(U_9 + V_9\right) - X\right)$$

$$\mathrm{X}\sin\left(\left(Z_9 T_9\right) - M_9 V_9 \sin\left(Z_9 T_9\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 =$ $1 - cos\, Z_9\ T_9$and $A_9 = U_9 - V_9$we obtain

$$4\, M_9^2 V_9^2 W_9^4 + 2\, M_9 V_9\ W_9^3[(M_9 - V_9)^2 + A_9\left(M_9 + V_9\right) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2\left[M_9 - V_9\right)^2\left(2\, A_9\ M_9 + Y_9 - X_9\right) + \left(2A_9\ M_9 - X_9\right)^2 - A_9\left(M_9 - V_9\right)\left(2A_9\ M_9 - X_9\right)$$

$$+ M_9 V_9\,(A_9^2 + 4\left(X_9 + Y_9\right))]$$

$$+ W_9\left[A_9^2 + 2X_9\right)\left(2\, A_9\ M_9 - X_9\right) + 2A_9 Y_9\,(M_9 + V_9) - M_9 A_9^3] + A_9^2\left(X_9\ M_9 - Y_9\right) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9(2, U_9 V_9) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting cos $(Z_9 T_9) = -1\ in\ (A3.3)$

We have applied the model to Kashmiri Shaivism. Whose elaborate rituals is enough to make a indivual raise from individual consciousness to cosmic consciousness which we have called Cosmic Consciousness.

It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) **Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP** The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic

Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** The group of unimodular transformations on the roots of polynomials over GF(2) is considered, and those polynomials with (e&eb) symmetries in (eb) the unimodular group are identified. The cross-correlation function between (e&eb) two maximum-length linear shift register sequences of (e) the same degree is shown to be (=) computable as an explicit linear transformation, in (eb) matrix form, on (e&eb) either one of the sequences, regarded as (=) a vector. The underlying vector space is (=) the "cyclotomic algebra," generated by (e) the cyclotomic cosets, or by what Gauss termed the **"periods" of the cyclotomic equation**. A variety of numerical examples are worked in detail. **Information Sciences Volume 1, Issue 1, December 1968, Pages 87-109 Theory of transformation groups of polynomials over GF(2) with applications to linear shift register sequences Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more link. Solomon W. Golomb** stop Authors study the P−V criticality and phase transition in (eb) the extended phase space of charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of (e) the system and its conjugate quantity is (=) the thermodynamic volume of the black hole. The black holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb) no P−V criticality and phase transition appear, while (e) for the black holes with a spherical horizon, even when the charge of the black hole is absent, t(eb) he P−V criticality and the small black hole/large black hole phase transition will appear, but it happens only in d=5 dimensions; when the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to have (e) the P−V criticality

and (e&eb)the small black hole/large black hole phase transition, there exists (ebb) an upper bound for the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the **Gauss-Bonnet coefficient** and Q is (=) the charge of the black hole. Authors calculate the critical exponents at the critical point and find that (eb) for all cases, (eb) they are the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09(2013)005 Cite as: arXiv:1306.6233 [gr-qc] (or arXiv:1306.6233v4 [gr-qc] for this version) **P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in IR but with (e&eb) improved UV behaviors. Here, authors study an IR modification which breaks ``softly'' (e) the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Authors obtain the black hole and cosmological solutions for **``arbitrary'' cosmological constant** that represent (eb) the analogs of the standard Schwarzschild-(A)dS solutions which can be asymptotically (A)dS as well as flat and authors discuss their thermodynamical properties. They also obtain solutions for (e) FRW metric with (e&eb) an arbitrary cosmological constant. Authors study its implication to (e&eb) the dark energy and find that it seems to be consistent with (e&eb) current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP.** The Einstein equations with (e&eb) a negative cosmological constant admit (e) black hole solutions which are asymptotic to (e) anti-de Sitter space. Like black holes in asymptotically flat space, these solutions have (e) thermodynamic properties including (e) a characteristic temperature and intrinsic entropy equal to (=) one quarter of the area of the event horizon in Planck units. There are however some important differences from (e) the asymptotically flat case. A black hole in anti-de Sitter space has (e) a minimum temperature which occurs when (e) its size is of the order of the characteristic radius of (e) the anti-de Sitter space. For larger black holes the red-shifted

temperature measured at (eb) infinity is greater. This means that such black holes have (e) positive specific heat and can be in stable equilibrium with (e&eb) thermal radiation at a fixed temperature. It also implies (eb) that the canonical ensemble exists for asymptotically anti-de Sitter space, unlike the case for asymptotically flat space. One can also consider the microcanonical ensemble. One can avoid (e) the problem that arises in asymptotically flat space of (e) having to put the system in (eb) a box with unphysical perfectly reflecting walls because (e) the gravitational potential of anti-de Sitter space acts as a box of finite volume. Communicated by A. Jaffe **Communications in Mathematical Physics December 1983, Volume 87, Issue 4, pp 577-588 Thermodynamics of black holes in anti-de Sitter space S. W. Hawking, Don N. Page STOP** In previous work authors introduced the concept of black-hole entropy, which they identified with (e&eb) the surface area of the black hole in question expressed in units of (e) the Planck length squared. Authors suggested that the appropriate generalization of the second law for a region containing a black hole is (=) that the black-hole entropy plus (+) the common entropy in the black-hole exterior never decreases. Here authors establish the validity of this law for the infall of an entropy-bearing system into (e&eb) a much larger and more massive generic stationary black hole. To do this authors determine (eb) a general lower bound for the increase in black-hole entropy, and an upper bound for the entropy of the system, while allowing for quantum effects at each stage. In passing authors show that the generalized second law is a statistical law which becomes (=) over-whelmingly probable in the limit of a macroscopic system. Authors also consider briefly more general situations. Finally, we give two simple examples of predictions made by the generalized second law for (e) black-hole formation processes. Received 17 September 1973 DOI:http://dx.doi.org/10.1103/PhysRevD.9.3292 **Generalized second law of thermodynamics in black-hole physics Jacob D. Bekenstein Phys. Rev. D 9, 3292 – Published 15 June 1974 STOP** The Lovelock gravity consists of (e) the dimensionally extended Euler densities. The geometry and horizon structure of (e) black hole solutions could be quite complicated in (eb) this gravity; however, authors find that some thermodynamic quantities

of the black holes like the mass, Hawking temperature and entropy have (e) simple forms expressed in terms of (e) horizon radius. The case with black hole horizon being a Ricci flat hypersurface is (=) particularly simple. In that case the black holes are (=) always thermodynamically stable with (e&eb) a positive heat capacity and their entropy still obeys (e&eb) the area formula, which is no longer valid for (e) black holes with positive or negative constant curvature horizon hypersurface. In addition, for black holes in (eb) the gravity theory of Ricci scalar plus a 2n-dimensional Euler density with (e&eb) a positive coefficient, thermodynamically stable small black holes always exist in (eb) D=2n+1 dimensions, which are absent in (eb) the case without (e) the Euler density term, while (e) the thermodynamic properties of the black hole solutions with (e&eb) the Euler density term are qualitatively similar to those of black holes without (e) the Euler density term as D>2n+1. Copyright © 2004 Published by Elsevier B.V. **Physics Letters B Volume 582, Issues 3–4, 4 March 2004, Pages 237–242 A note on thermodynamics of black holes in Lovelock gravity Rong-Gen CAI STOP** There are a number of similarities between black-hole physics and (e, eb, =) thermodynamics. Most striking is the similarity in the behaviors of black-hole area and (e&eb) of entropy: Both quantities tend to increase (eb+) irreversibly. In this paper we make this similarity the basis of (e) a thermodynamic approach to black-hole physics. After a brief review of the elements of the theory of information, authors discuss black-hole physics from the point of view of information theory. Authors show that it is natural to introduce (eb) the concept of black-hole entropy as (=) the measure of information about (e&eb) a black-hole interior which is (=) inaccessible to an exterior observer. Considerations of simplicity and consistency, and dimensional arguments indicate (eb) that the black-hole entropy is equal to (=) the ratio of the black-hole area to the square of the Planck length times a dimensionless constant of (e) order unity. A different approach making use of the specific properties of Kerr black holes and of concepts from information theory leads to (eb) the same conclusion, and suggests (eb) a definite value for the constant. The physical content of (e) the concept of black-hole entropy derives from (e) the following generalized version of the second law: When common entropy goes down a

black hole, (eb) the common entropy in (eb) the black-hole exterior plus the black-hole entropy never decreases. The validity of this version of the second law is supported by (e) an argument from information theory as well as by several examples. Received 2 November 1972 DOI:http://dx.doi.org/10.1103/PhysRevD.7.2333 **Black Holes and Entropy Jacob D. Bekenstein Phys. Rev. D 7, 2333 – Published 15 April 1973 STOP** The effects of power-law Maxwell field on the Van der Waals like phase transition of higher-dimensional dilaton black holes are probed in detail. It is shown that the Smarr relation gains corrections due to (e) the effects of (e) both the power-law Maxwell field and (e&eb) the dilaton field while (e) thermodynamic volume is exactly the same as (=) that of Einstein-Maxwell-dilaton black holes. Authors successfully derive the analytic solutions of (e) critical point and carry out (e&eb) some check to ensure that these critical quantities are (=) positive. It is shown that the constraint on the parameters turns out to be (=) 0<\alpha^2<1, which is tighter than that in (eb) the non-extended phase space. It is also shown (eb) that these critical quantities and (e&eb) the ratio P_cv_c/T_c are affected by (e&eb) the power-law Maxwell field. Moreover, critical exponents are found to coincide with those of other AdS black holes, showing (eb) the powerful influence of mean field theory. Comments: references added Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1601.05500 [gr-qc] (or arXiv:1601.05500v2 [gr-qc] for this version **Effects of power-law Maxwell field on the Van der Waals like phase transition of higher dimensional dilaton black holes Jie-Xiong Mo, Gu-Qiang Li, Xiao-Bao Xu STOP** A consistent and unified picture for (e) critical phenomena of charged AdS black holes in f(R) gravity is drawn in this paper. Firstly, we investigate the phase transition in canonical ensemble. Authors derive the explicit solutions corresponding to (e&eb) the divergence of CQ. The two solutions merge into (e&eb) one when (e) the condition Qc=−13R0−−−√ is satisfied. The curve of specific heat for Q<Qc has two divergent points and can be divided into (e&eb) three regions. Both the large radius region and the small radius region are (=) thermodynamically stable with (e&eb) positive specific heat while (e) the medium radius region is unstable with (e&eb) negative specific heat. However, when Q>Qc, the specific heat is

always positive, implying (eb) the black holes are locally stable and no phase transition will take place. Secondly, both the T−r+ curve and T−S curve f(R) AdS black holes are investigated and they exhibit Van der Vaals like behavior as the P−v curve in the former research. Critical physical quantities are obtained and they are consistent with (=) those derived from the specific heat analysis. Authors carry out numerical check of Maxwell equal area law for the cases Q=0.2Qc, 0.4Qc, 0.6Qc, 0.8Qc. The relative errors are (=) amazingly small and can be negligible. So the Maxwell equal area law holds for (e) T−S curve of f(R) black holes. Thirdly, authors establish geometrothermodynamics for (e) f(R) AdS black hole to examine the phase structure. It is shown that the Legendre invariant scalar curvature R would diverge exactly where (e) the specific heat diverges. To summarize, the above three perspectives are consistent with (e&eb) each other, thus providing (eb) a unified picture which deepens the understanding of (e) critical phenomena of (e) f(R) AdS black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1602.01251 [gr-qc] (or arXiv:1602.01251v1 [gr-qc] for this version) **A consistent and unified picture for critical phenomena of f(R) AdS black holes Jie-Xiong Mo, Gu-Qiang Li, Yu-Cheng Wu STOP** The state equation of the charged AdS black hole is reviewed in (e&eb) the T−r plane. Thinking of the phase transition, the T−S, P−V, P−ν graphs are plotted and then the equal area law is used (e) in the three cases to get (eb) the phase transition point (P, T). The analytical phase transition point relations for (e) P-T of charged AdS black hole has been obtained successfully. By comparing the three results, authors find that the equal area law possibly cannot (e) be used directly for P−ν plane. According to the T−S, P−V results, they plot (e&eb) the P−T−Q graph and find that for a highly charged black hole a very low temperature condition is required for (e) the phase transition. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI:10.1140/epjc/s10052-015-3641-0 Cite as: arXiv: 1503.07658 [gr-qc] (or arXiv:1503.07658v1 [gr-qc] for this version) **A note on Maxwell's equal area law for black hole phase transitionShan-Quan Lan, Jie-Xiong Mo, Wen-Biao Liu STOP** Recently, extended phase space thermodynamics of Lovelock AdS black holes has been of great interest. To provide

insight from (e) a different perspective and gain (e) a unified phase transition picture, non-extended phase space thermodynamics of (e) (n+1)-dimensional charged topological Lovelock AdS black holes is (=) investigated detailedly in (eb) the grand canonical ensemble. Specifically, the specific heat at (eb) constant electric potential is calculated and phase transition in the grand canonical ensemble is discussed. To probe the impact of the various parameters, they utilize (e) the control variate method and solve (e) the phase transition condition equation numerically for the case k=1, −1. There are two critical points for (e) the case n=6, k=1 while there is only one for other cases. For k=0, there exists no phase transition point. To figure out the nature of phase transition in the grand canonical ensemble, authors carry out an analytic check of the analog form of (e) Ehrenfest equations proposed by Banerjee et al. It is shown that Lovelock AdS black holes in (eb) the grand canonical ensemble undergo (e&eb) a second order phase transition. To examine the phase structure in (eb) the grand canonical ensemble, authors utilize the thermodynamic geometry method and calculate both the Weinhold metric and Ruppeiner metric. It is shown that for both analytic and graphical results that the divergence structure of (e) the Ruppeiner scalar curvature coincides with (=) that of the specific heat. Research provides (eb) one more example that Ruppeiner metric serves as (=) a wonderful tool to (e) probe the phase structures of (e) black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C75 (2015) 5, 211 DOI:10.1140/epjc/s10052-015-3434-5 Cite as: arXiv: 1503.01956 [gr-qc] (or arXiv: 1503.01956v1 [gr-qc] for this version **Non-extended phase space thermodynamics of Lovelock AdS black holes in grand canonical ensemble Jie-Xiong Mo, Wen-Biao Liu STOP** To provide an analytic verification of the nature of phase transition at the critical point of P−V criticality, the original expressions of Ehrenfest equations have been introduced directly. By treating the cosmological constant and (e&eb) its conjugate quantity as (=) thermodynamic pressure and volume respectively, authors carry out analytical check of (e) classical Ehrenfest equations. To show that approach is universal, authors investigate not only higher-dimensional charged AdS black holes (e&eb), but also rotating AdS black holes. Not only are the examples

of Einstein gravity shown, but also the example of modified gravity is presented for (e) Gauss-Bonnet AdS black holes. The specific heat at constant pressure CP, the volume expansion coefficient α and the isothermal compressibility coefficient κT are found to diverge exactly at the critical point. It has been verified that both Ehrenfest equations hold at (eb) the critical point of P−V criticality in (eb) the extended phase spaces of (e) AdS black holes. So the nature of the critical point of P−V criticality of AdS black holes has been demonstrated analytically to be (=) a second-order phase transition. These results are consistent with (e&eb) the nature of liquid-gas phase transition at the critical point. In this sense, extant research would deepen the understanding of (e) the relations of AdS black holes and (e&eb) liquid-gas systems. Moreover, successful approaches to introduce (eb) the original expressions of Erhenfest equations directly into (e&eb) black hole phase transition research demonstrate again (eb) that black hole thermodynamics is closely related (e&eb) to classical thermodynamics, which allows (eb) us to borrow techniques from (e) classical thermodynamics to investigate the thermodynamics of black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 89, 084057 (2014) DOI: 10.1103/PhysRevD.89.084057 Cite as:arXiv: 1404.3872 [gr-qc] (or arXiv:1404.3872v1 [gr-qc] for this version) **Ehrenfest scheme for P−V criticality of higher dimensional charged black holes, rotating black holes and Gauss-Bonnet AdS black holes Jie-Xiong Mo, Wen-Biao Liu STOP** Aiming at a unified phase transition picture of (e) the charged topological black hole in (eb) Ho\v{r}ava-Lifshitz gravity, authors investigate this issue not only in (eb) canonical ensemble with (e&eb) the fixed charge case but also in grand-canonical ensemble with (e&eb) the fixed potential case. Authors firstly perform the standard analysis of the specific heat,(e&eb) the free energy and the Gibbs potential, and (e&eb) then study its geometrothermodynamics. It is shown that the local phase transition points (eb) not only witness the divergence of (e) the specific heat, but also witness the minimum temperature and the maximum free energy or Gibbs potential. They also witness the divergence of (e) the corresponding thermodynamic scalar curvature. No matter which ensemble is chosen, the metric constructed can

successfully produce (eb) the behavior of the thermodynamic interaction and (e&eb) phase transition structure while other metrics failed to (e) predict the phase transition point of (e) the charged topological black hole in former literature. In grand-canonical ensemble, authors have discovered the phase transition which has not been reported before. It is similar to (e) the canonical ensemble in which the phase transition only takes place when k=−1. But it also has its unique characteristics that the location of the phase transition point depends on (eb) the value of potential, which is different from (e) the canonical ensemble where (e) the phase transition point is (=) independent of the parameters. After an analytical check of Ehrenfest scheme, authors find that the new phase transition is (=) a second order one. It is also found that the thermodynamics of (e) the black hole in Horava-Lifshitz gravity is quite different from that in (eb) Einstein gravity. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP1310(2013)056 DOI:10.1007/JHEP10(2013)056 Cite as: arXiv:1404.2497 [gr-qc] (or arXiv:1404.2497v1 [gr-qc] for this version) **A unified phase transition picture of the charged topological black hole in Horava-Lifshitz gravity Jie-Xiong Mo, Xiao-Xiong Zeng, Gu-Qiang Li, Xin Jiang, Wen-Biao Liu STOP** To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case $\beta\to\infty$ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness condition of (e) entropy. It is shown

that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure authors have discovered. It is shown that Legendre invariant thermodynamic scalar

curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws, leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions, geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged AdS black holes by considering the case of the RN-AdS black holes

surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase transition. For the case $\omega_q=-2/3$ authors derive analytic expressions of (e) critical physical quantities, while (e) for cases $\omega_q\neq-2/3$ appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this

version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=) chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy

Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed

quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions, while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma?**

All this consists of the act of consciousness alone", says Vasubandhu, "because unreal entities appear, just as a man with defective vision sees unreal hair or a moon, etc."He points to dreams as examples of purely subjective constructs which appear to be objective realities. The apparent reality dreams possess is not derived from any concrete, objective world, but merely from the idea of objectivity. While the Yogacara does not say that an idea has, for example, spatial attributes, it does have a form manifesting them. While he agrees with the Saiva idealist that

appearances have no independent existence apart from their appearing to consciousness, he maintains that for this reason they are unreal. The creativity of consciousness consists in its diversification in many modes having apparent externality; it is not a creation of objects. While the Kashmiri Saivite agrees that the world is pure conscious- ness alone, he maintains that it is such because it is a real creation of consciousness. The effect is essentially identical with the cause and shares in its reality. Matter and the entire universe are absolutely real, as 'congealed' (sty ana) or 'contracted' (samkucita) forms of consciousness. "This God of consciousness", writes Ksemaraja, "generates the universe and its form is a condensation of His own essence (rasa)" m By boiling sugarcane juice it condenses to form treacle, brown sugar and candy which retains its sweetness. Similarly, consciousness abides unchanged even though it assumes the concrete material form of the five gross elements. The same reality thus abides equally in gross and subtle forms. Consequently no object is totally insentient. Even stones bear a trace (vasana) of consciousness, although it is not clearly apparent because it is not associated with the vital breath (prana) and other components of a psycho-physical organism. Somananda goes so far as to affirm that physical objects, far from being insentient, can only exist insofar as they are aware of themselves as existing. The jar performs its function because it knows itself to be its agent. Indeed, all things are pervaded by consciousness and at one with it and hence share in its omniscience. Thus, Siva, Who perceives Himself in the form of physical objects, is the one ultimate reality. "The jar knows because it is of my nature", writes Somananda, "and I know it because I am of the jar's nature. I know because I am of Sadasiva's nature and He knows because He is of my nature; Yajnadatta [knows] because he is of Siva's nature and Siva [knows] because He is of Yajnadatta's nature". Integral Monism of Kashmiri Saivism Everything in this sense is directly perceived by absolute conscious- ness, and this direct perception (pratyaksa) unifies the knowable into a single, undivided whole. This is the central concept behind a doctrine originally expounded by Narasimha called 'the non-dualism of direct perception' (pratyaksadvaita). This states that consciousness is essentially perceptive and that its perception of all things operates throughout the universe. Insofar

as phenomena are clearly evident (sphufa) to us, everything is directly perceived by absolute consciousness, with which our individual consciousness is identical. This direct perception unfolds everywhere; the one true reality, it is alone and without companion or rival (nihsapatna). Even though it remains one, it can, by its very nature, perceive distinctions (bheda) between one entity and another, without this engendering any division within it. We distinguish between two entities in empirical terms on the basis of their mutual exclusion (anyonyabhava). The relative distinction {bheda) between them is essentially the perceived difference between their respective characteristics. Despite this difference they are united within the purview of a single cognition insofar as they are equally both manifest appearances. This cognition is the undivided essence (rasa) or 'own nature' (svabhava) of both. Encompassed by the 'fire of consciousness', there is no essential difference between them. Just as when an emerald and ruby reflect each other's light, the ruby is reddish-green and the emerald greenish-red, similarly everything is connected with everything else as part of the single variegated (vicitra) cognition of absolute consciousness. Mahesvarananda writes: The Supreme Lord's unique state of emotivity (asadharanabhava) is the outpouring of pure Being (mahasatta). It is manifest as the brilliance (sphuratta) of the universe which, if we ponder deeply, [is realized to be] the single flavour (ekarasa) of the essence of Beauty which is the vibration of the bliss of one's own nature. In this way all things are in reality one although divided from the one another sharing as they do the 'single flavour' (ekarasa) of the pure vibration of consciousness. Kashmiri Saiva Realism Kashmiri Saivism as a whole has been variously called a form of 'realistic idealism', 'monistic idealism', ‘idealistic monism' and 'concrete monism'. It is easy to understand why Kashmiri Saivism is The Doctrine of Vibration said to be 'idealistic' and 'monistic', but in what sense is it also 'realistic*? The answer to this question is of no small importance in trying to under- stand the central idea behind its metaphysics and the fundamental importance of the concept of Spanda, in this seemingly impossible marriage between monistic idealism and pluralistic realism. The Kashmiri Saiva approach understands the world to be a symbol of the absolute, that is, as the manner in which it presents itself to us. Again we can

contrast this view with that of the Advaita Vedanta. The Advaita Vedanta understands the world to be an expression of the absolute insofar as it exists by virtue of the absolute's Being. Being is understood to be the real unity which underlies empirically manifest separateness and as such is never empirically manifest. It is only transcendentally actual as 'being-in-itself. The Kashmiri Saiva position represents, in a sense, a reversal of this point of view. The nature of the absolute, and also that of Being, is onceived as an eternal becoming (satatodita), a dynamic flux or Spanda, 'the agency of the act of being'. It is dentified with the concrete actuality of the fact of appearing, not passive unmanifest Being. Appearance (abhasa) alone is real.ppearing (prakasamanatva) is equivalent to the fact of being (astitva). Ksemaraja writes in his commentary on the Stanzas on Vibration: Indeed, all things are manifest because they are nothing but manifestation. The point being that nothing is manifest apart from manifestation The absolutely unmanifest, from this point of view, can have as little existence as the space in a lattice window of a sky-palace. Nay, even less, because even that space can appear as an imagined mage manifest within consciousness. Everything is real according to the manner in which it appears. Even an illusion is in this sense real, insofar as it appears and is known in the manner in which it appears. The empirical and the real are identical categories of thought. As Abhinava says: Thus this is the supreme doctrine (upanisatf), namely that, when- ever and in whatever form [an entity] appears, that then is its particular nature. Perhaps at this stage a brief comparison with Heidegger's ideas might prove to be enlightening and not altogether out of place. According to Heidegger's phenomenology of Being, reality is intelligible in a two-fold Integral Monism of Kashmiri Saivism manner as 'phenomenon' and 'logos'. Heidegger defines what he means by 'phenomenon' as: "that-which-shows-itself. The manifest . . . phenomena are then the collection of that which lies open in broad daylight or can be brought to the light of day — what the Greeks at times implicitly identified as 'ta onta' (the things-which-are)".

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb) heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna**

**Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version) **Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb) strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel

period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar (=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI: 10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills (SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear

differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black**

**holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-

relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb)

the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this

version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstate thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11 August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of

entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.** Characterizing the behaviour of strongly coupled

quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at λ=λc, the correlation length diverges as ξ∼|λ−λc|−ν, while (e) the gap between the ground state and first excited state vanishes as (e) ΔE∼ξ−z∼|λ−λc|zν. Exactly at the critical point λ=λc, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as ω∼|k−k0|z. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-**

**equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004 STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element.

**Model has been applied for the following results:**

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| **Searches for Sterile Neutrinos with the IceCube Detector M. G. Aartsen et al. (IceCube Collaboration) Phys.**<br><br>**Rev. Lett 117, 071801 – Published 8 August 2016 Physics logo See Viewpoint: Hunting the Sterile Neutrino** |
| **(1) The IceCube neutrino telescope at the South Pole has measured the atmospheric muon neutrino spectrum as (=) a function of zenith angle and (e&eb) energy in the approximate 320 GeV to 20 TeV range, to search for (e) the oscillation signatures of (e) light sterile neutrinos.** |
| **(2) No evidence for anomalous νμ or ν¯μ disappearance is observed in (eb) either of two independently developed analyses, each using one year of (e) atmospheric neutrino data.** |
| **(3) New exclusion limits are placed on (e&eb) the parameter space of (e) the 3+1 model, in which muon antineutrinos experience (e) a strong Mikheyev-Smirnov-Wolfenstein-resonant oscillation.** |
| **(4) The exclusion limits extend to (e) sin22024≤0.02 at Δm2∼0.3 eV2 at the 90% confidence level.** |
| **(5) The allowed region from (e) global analysis of appearance experiments, including (e) LSND and MiniBooNE, is excluded at approximately the 99% confidence level for (e) the global best-fit value of \|Ue4\|2. Received 5 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.117.071801 © 2016 American Physical Society** |
| **Synopsis: Searching for Majorana Neutrinos August 16, 2016The KamLAND-Zen collaboration has run the most sensitive search to date for a radioactive decay that could reveal whether neutrinos are Majorana fermions. Synopsis figureKamLAND-Zen Collaboration** |
| **(6) Are neutrinos and (=) antineutrinos the same particles, as hypothesized by (e) physicist Ettore Majorana?** |
| **(7) If true, this "Majorana" nature of neutrinos could explain (eb) why neutrinos have (e) mass and why the Universe has (e) more matter than antimatter.** |
| **(8) Using an underground neutrino detector in Japan, the KamLAND-Zen collaboration has carried out the most sensitive search to date for a process called neutrinoless double-beta decay, which could confirm (eb) Majorana's hypothesis.** |
| **(9) The results set stringent constraints on the possible properties of (e) Majorana neutrinos.** |
| **(10) Double-beta decay is (=) a nuclear process in which two neutrons turn into (e&eb) two protons (or vice versa), emitting (eb) two antineutrinos.** |
| **(11) If neutrinos are Majorana fermions, (eb) the two antineutrinos could annihilate each other, and some double-beta decays would emit (eb) no antineutrinos.** |
| **(12) However, neutrinoless decay is believed to be (=) extremely rare. Searches of the process using (e) large amounts of the 35 naturally occurring (eb) isotopes that can undergo (e&eb) double-beta decay (such as xenon-136) have so far come up empty.** |

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| **(13) The KamLAND-Zen collaboration boosted their experimental sensitivity to (e) neutrinoless double-beta decays through (e&eb) a combination of factors.** |
| **(14) The researchers deployed clean detectors with (e&eb) very low background noise.** |
| **(15) And they used an unprecedented amount of xenon-136, which they purified to (e) remove any radioactive contaminants that could produce (eb) unwanted signals in the detectors.** |
| **(16) Combining data from several years, they improved (eb) the limit on the probability of neutrinoless double-beta decay by (e) sixfold compared to previous searches.** |
| **(17) This probability corresponds to (e&eb) a decay lifetime whose staggering value exceeds (e) 1026 years.** |
| **(18) This measurement allowed (eb) the researchers to set (eb) the most stringent upper limit on (e&eb) the possible mass of Majorana neutrinos (the particle must be lighter than 61–165 meV). This research is published in Physical Review Letters. –Matteo Rini Matteo Rini is the Deputy Editor of Physics.** |
| **How to Reach a Thousand-Second in-Plane Polarization Lifetime with 0.97−GeV/c Deuterons in a Storage Ring G. Guidoboni et al. (JEDI Collaboration) Phys. Rev. Lett. 117, 054801 – Published 28 July 2016 Physics logo See Synopsis: Deuterons Spin together for a Thousand Seconds** |
| **(19) Authors observe a deuteron beam polarization lifetime near 1000 s in (eb) the horizontal plane of (e) a magnetic storage ring (COSY).** |
| **(20) This long spin coherence time is maintained through (e&eb) a combination of beam bunching, electron cooling, (e&eb) sextupole field corrections, and (e&eb) the suppression of collective effects through (e&eb) beam current limits.** |
| **(21) This record lifetime is required for (e) a storage ring search for (e) an intrinsic electric dipole moment on (e&eb) the deuteron at a statistical sensitivity level approaching 10−29 e cm.** |
| **Synopsis: Deuterons Spin together for a Thousand Seconds July 28, 2016 Researchers set a new record for the in-plane spin-alignment lifetime of deuterons circulating in a magnetic storage ring. Synopsis figure E. Stephenson/Indiana University Center for Spacetime Symmetries Sime sentences and figures are deleted due to spatial restrictions.** |
| **(22) The electric dipole moment (EDM) of (e) a fundamental particle or a simple nucleus is believed to be (=) extremely small.** |
| **(23) One might therefore think that it isn't worth measuring this quantity. But it is. A nonzero EDM would signal (eb) a broken symmetry of (e) nature known as CP violation, which is required to (e) explain why the Universe has (e) more matter than antimatter.** |
| **(24) Edward Stephenson from Indiana University, Bloomington, and colleagues has now made a significant stride towards determining (eb) the EDM of the deuteron using (e) a deuteron beam circulating in (e&eb) a magnetic storage ring.** |

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| **(25) They have been able to preserve the alignment of the deuterons' spins for (e) a chunk of time that may be long enough to determine (eb) the deuteron's EDM.** |
| **(26) In experiments of this kind, the EDM can be determined by (e) measuring the rotation of (e) the initially aligned spins of (e) the beam particles as they travel around (e&eb) the ring.** |
| **(27) To obtain a measurement sufficiently accurate to explore (e&eb) CP violation predicted by (e) theories beyond the standard model, the spin alignment has to be (=) preserved for about 1000 s.** |
| **(28) But changes in particle speed—associated with (e&eb) position or direction variations among (e&eb) the particles in the beam—lead to (eb) a rapid loss of the alignment in the plane of the ring.** |
| **(29) Stephenson and co-workers have managed to retain (eb)the alignment of deuterons for the required 1000 s, which is 3 orders of magnitude longer (eb) than the previous record (obtained for an electron beam).** |
| **(30) Their success resulted from (e) capturing the beam into a tight bunch to reduce (e) position and direction spread and from (e) shaping the ring's magnetic field to(e) ensure all of the deuteron orbits around (e&eb) the ring have the same length, despite (e) direction variations. This research is published in Physical Review Letters. –Ana Lopes Ana Lopes is a Senior Editor of Physics.** |
| **Cotunneling Drag Effect in Coulomb-Coupled Quantum Dots A. J. Keller, J. S. Lim, David Sánchez, Rosa López, S. Amasha, J. A. Katine, Hadas Shtrikman, and D. Goldhaber-Gordon Phys. Rev. Lett. 117, 066602 – Published 4 August 2016 Physics logo See Synopsis: Coulomb Drag in a Double Dot** |
| **(31) In Coulomb drag, a current flowing in one conductor can induce (eb) a voltage across an adjacent conductor via (e&eb) the Coulomb interaction.** |
| **(32) The mechanisms yielding drag effects are not always understood, even though drag effects are (=) sufficiently general to be (=) seen in many low-dimensional systems.** |
| **(33) In this Letter, authors observe Coulomb drag in (eb) a Coulomb-coupled double quantum dot and, through (e&eb) both experimental and theoretical arguments, identify cotunneling as (=) essential to obtaining a correct qualitative understanding of the drag behavior.** |
| **Synopsis: Coulomb Drag in a Double Dot August 4, 2016 Electric current passing through a quantum dot can generate current in a nearby dot through a coordinated tunneling mechanism. Synopsis figureA. J. Keller et al., Phys. Rev. Lett. (2016)** |
| **(34) Electrons moving through (e&eb) a conductor can induce (eb) the movement of electrons of a nearby, but isolated, conductor.** |
| **(35) This so-called Coulomb drag is well understood in 2D layered materials, but the theory for nanowires and quantum dots does not (e) fully explain experiments.** |
| **(36) A new study of Coulomb drag in two nearby quantum dots shows (eb) that simultaneous tunneling of electrons—called cotunneling—can help explain the low-temperature behavior of (e) the system.** |

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| **(37) Electric current flows through (e&eb) a quantum dot by the tunneling of (e) electrons.** |
| **(38) When two dots are placed next to each other, the active, or lead, dot can drag current through (e&eb) the passive dot by (e) a process called sequential tunneling.** |
| **(39) In this case, an electron first tunnels onto (e&eb) the active dot, increasing (eb) the probability of tunneling onto (e&eb) the passive dot.** |
| **(40) Thus the passive dot follows (e&eb) the active dot's lead, but always one step behind.** |
| **(41) Sequential tunneling explains (eb) Coulomb drag observations at (eb) high temperature, but no studies have investigated cold temperatures where (e) quantum coherence effects may play a role.** |
| **(42) Andrew Keller from the California Institute of Technology, Pasadena, and his colleagues explored Coulomb drag at (eb) sub-Kelvin temperatures in a pair of quantum dots coupled by (e) capacitance.** |
| **(43) The team applied a current-driving voltage across (e&eb) one dot and detected a drag current in (eb) the other dot, even in regimes where (e) the driving voltage was very small.** |
| **(44) This conflicts with (e&eb) the sequential tunneling model, which predicts (eb) a threshold driving voltage below which (e) the drag current should be zero.** |
| **(45) The team showed that the inclusion of cotunneling into (eb) the drag theory explains the observations.** |
| **(46) Cotunneling could be important to efforts aimed at harvesting a unidirectional current from (e) a fluctuating voltage supply. This research is published in Physical Review Letters. –Michael Schirber Michael Schirber is a Corresponding Editor for Physics based in Lyon, France.** |
| **Maybe it wasn't the Higgs particle after all November 7, 2014 Maybe it wasn't the Higgs particle after all Associate Professor Mads Toudal Frandsen, University of Southern Denmark Credit: University of Southern Denmark** |
| **(47) Last year CERN announced the finding of a new elementary particle, the Higgs particle. But maybe it wasn't the Higgs particle, maybe it just looks like it. And maybe it is not alone. Many calculations indicate that the particle discovered last year in the CERN particle accelerator was (=) indeed the famous Higgs particle.** |
| **(48) Physicists agree that the CERN experiments did find a new particle that had never been seen before, but according to (e&eb) an international research team, there is (=) no conclusive evidence that the particle was indeed the Higgs particle.The research team has scrutinized the existing scientific data from CERN about the newfound particle and published their analysis in the journal Physical Review D. A member of this team is Mads Toudal Frandsen, associate professor at the Center for Cosmology and Particle Physics Phenomenology, Department of Physics, Chemistry and Pharmacy at the University of Southern Denmark. "The CERN data is generally taken as (=) evidence that the particle is the Higgs particle.** |

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| **(49) It is true that the Higgs particle can explain (eb) the data but there can be other explanations, we would also get this data from other particles", Mads Toudal Frandsen explains.** |
| **(50) The researchers' analysis does not (e) debunk the possibility that CERN has discovered the Higgs particle. That is still possible - but it is equally possible that it is a different kind of particle. "The current data is not precise enough to determine exactly what the particle is. It could be a number of other known particles", says Mads Toudal Frandsen. But if it wasn't the Higgs particle, which was found in CERN's particle accelerator, then what was it? "We believe that it may be a so-called techni-higgs particle. This particle is in some ways similar to the Higgs particle - hence half of the name", says Mads Toudal Frandsen.** |
| **(51) Although the techni-higgs particle and Higgs particle can easily be confused in experiments, they are two very different particles belonging to (e) two very different theories of how the universe was created.** |
| **(52) The Higgs particle is the missing piece in (eb) the theory called the Standard Model.** |
| **(53) This theory describes (eb) three of the four forces of nature. But it does not (e) explain what dark matter is - the substance that makes up (eb) most of the universe. A techni-higgs particle, if it exists, is (=) a completely different thing:** |
| **(54) "A techni-higgs particle is not (e) an elementary particle.** |
| **(55) Instead, it consists of so-called techni-quarks, which we believe are elementary. Techni-quarks may bind together in various ways to form (eb) for instance techni-higgs particles, while other combinations may form (eb) dark matter. Authors therefore expect to find several different particles at the LHC, all built by (e) techni-quarks", says Mads Toudal Frandsen.** |
| **(56) If techni-quarks exist, (eb) there must be a force to bind them together so that (e) they can form particles.** |
| **(57) None of the four known forces of nature (gravity, the electromagnetic force, the weak nuclear force and the strong nuclear force) are (=) any good at binding techni-quarks together.** |
| **(58) There must therefore be (=) a yet undiscovered force of nature. This force is called the the Technicolor force.** |
| **(59) What was found last year in CERN's accelerator could thus be (=) either the Higgs particle of the Standard Model or a light techni-higgs particle, composed of (e) two techni-quarks.** |
| **(60) Mads Toudal Frandsen believes that more data from CERN will probably be able to determine (eb) if it was a Higgs or a techni-higgs particle. If CERN gets an even more powerful accelerator, it will in principle be able to observe techni-quarks directly. More information: Technicolor Higgs boson in the light of LHC data. Phys. Rev. D 90, 035012th Alexander Belyaev, Matthew S. Brown, Roshan Foadi, and Mads T. Frandsen. journals.aps.org/prd/abstract/10.1103/PhysRevD.90.035012. On Arxiv: arxiv.org/abs/1309.2097 Journal reference: Physical Review D search and more info website: Provided by: University of Southern Denmark search and more info website** |

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| **The inflationary brane-antibrane universeClifford P. Burgess1,2, Mahbub Majumdar3, Detlef Nolte1, Fernando Quevedo3, Govindan Rajesh1 and Ren-Jie Zhang1 Published 13 August 2001 • Journal of High Energy Physics, Volume 2001, JHEP07(2001)** |
| **(61) Authors show how the motion through (e&eb) the extra dimensions of a gas of branes and antibranes can, under certain circumstances, produce (eb) an era of inflation as seen by (e) observers trapped on (e&eb) a 3-brane, with (e&eb) the inflaton being the inter-brane separation.** |
| **(62) Although most of discussion refers to (e) arbitrary ps, when we need to be specific they assume that they are (=) D-branes of type-II or type-I string theory.** |
| **(63) For realistic brane couplings, such as those arising in (eb) string theory, the inter-brane potentials are (=) too steep to inflate (eb) the universe for acceptably long times.** |
| **(64) However, for special regions of the parameter space of (e) brane-antibrane positions the brane motion is (=) slow enough for there to be (=) sufficient inflation.** |
| **(65) Inflation would be more generic in models where (e) the inter-brane interactions are much weaker.** |
| **(66) The spectrum of primordial density fluctuations predicted has (e) index n slightly less than 1, and acceptable amplitude, provided (e) that the extra dimensions have (e) linear size 1/r ~ 1012 GeV.** |
| **(67) Reheating occurs as in hybrid inflation, with (e&eb) the tachyonic instability of (e) the brane-antibrane system taking over for (e) small separations.** |
| **(68) The tachyon field can induce (eb) a cascade mechanism within which higher-dimension branes annihilate (e) into lower-dimension ones.** |
| **(69) They argue that such a cascade naturally stops with (e&eb) the production of 3-branes in (eb) 10-dimensional string theory.** |
| **Effect of the neutron lifetime on processes in the early universe A. V. Chechkin, A. V. Ivanchik, A. P. Serebrov, S. V. Bobashev First Online: 27 July 2016 DOI: 10.1134/S1063784216070069 Chechkin, A.V., Ivanchik, A.V., Serebrov, A.P. et al. Tech. Phys. (2016) 61: 1101. Doi: 10.1134/S1063784216070069** |
| **(70) The influence of the neutron lifetime on (e&eb) the abundance of light elements produced during (e&eb) the primordial nucleosynthesis early in (eb) the birth of the Universe is considered.** |
| **(71) Among light elements, namely, D, 3He, 4He, and 7Li, 4He proves to be (=) most sensitive to neutron lifetime τn.** |
| **(72) Astronomic data on the light element abundance also provide (eb) the best accuracy for 4He.** |
| **(73) The solution of a number of problems discussed in this paper requires (e) improving the accuracy of observations for the 4He abundance and refining the value of τn.** |

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| **Application of a Theory and Simulation based Convective Boundary Mixing model for AGB Star Evolution and Nucleosynthesis U. Battino, M. Pignatari, C. Ritter, F. Herwig, P. Denisenkov, J.W. Den Hartogh, R. Trappitsch, R. Hirschi, B. Freytag, F. Thielemann, B. Paxton (Submitted on 19 May 2016)** |
| **(74) The s-process nucleosynthesis in Asymptotic Giant Branch (AGB) stars depends on (e&eb) the modeling of (e) convective boundaries.** |
| **(75) Authors' present models and s-process simulations that adopt (e&eb) a treatment of convective boundaries based on (e) the results of hydrodynamic simulations and on (e&eb) the theory of mixing due to (e) gravity waves in (eb) the vicinity of convective boundaries.** |
| **(76) Hydrodynamics simulations suggest (eb) the presence of convective boundary mixing (CBM) at (eb) the bottom of the thermal pulse-driven convective zone.** |
| **(77) Similarly, convection-induced mixing processes are proposed for (e) the mixing below the convective envelope during (e&eb) third dredge-up where (e) the 13C pocket for (e) the s process in AGB stars forms.** |
| **(78) In this work apply apply a CBM model motivated by (e) simulations and theory to (e&eb) models with initial mass M = 2 and M = 3M?, and with (e&eb) initial metal content Z = 0:01 and Z = 0:02.** |
| **(79) As reported previously, the He-intershell abundance of 12C and 16O are increased by (e) CBM at the bottom of (e) pulse-driven convection zone.** |
| **(80) This mixing is affecting (e&eb) the 22Ne (alpha, n) 25Mg activation and (e&eb) the s-process effciency in the 13C-pocket.** |
| **(81) In our model CBM at the bottom of the convective envelope during (e&eb) the third dredgeup represents (eb) gravity wave mixing.** |
| **(82) Authors take further into account that hydrodynamic simulations indicate (eb) a declining mixing efficiency already about (e&eb) a pressure scale height from (e) the convective boundaries, compared to (e&eb) mixing-length theory.** |
| **(83) The final s-process abundances are characterized by (e) 0:36 < [s=Fe] < 0:78 and the heavy-to-light s-process ratio is (=) 0:23 < [hs=ls] < 0:45.** |
| **(84) Finally, authors compare results with (e&eb) stellar observations, pre-solar grain measurements and previous work. Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI:10.3847/0004-637X/827/1/30 Cite as: arXiv:1605.06159 [astro-ph.SR] (or arXiv:1605.06159v1 [astro-ph.SR] for this version)** |
| **MESA and NuGrid simulations of classical novae: CO and ONe nova nucleosynthesis Pavel A. Denissenkov (UVic), James W. Truran (Chicago), Marco Pignatari (Basel), Reto Trappitsch (Chicago), Christian Ritter (Frankfurt), Falk Herwig (UVic), Umberto Battino (Basel), Kiana Setoodehnia (McMaster), B. Paxton (UC Santa Barbara)** |

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| **(85) Classical novae are the result of (e) thermonuclear flashes of hydrogen accreted by (e) CO or ONe white dwarfs, leading (eb) eventually to the dynamic ejection of (e) the surface layers.** |
| **(86) These are observationally known to be (=) enriched in heavy elements, such as (=) C, O and Ne that must originate in (eb) layers below the H-flash convection zone.** |
| **(87) Building on their previous work, they now present stellar evolution simulations of (e) ONe novae and provide (eb) a comprehensive comparison of models with (e&eb) published ones.** |
| **(88) Some of authors models include exponential convective boundary mixing to account for (e) the observed enrichment of (e) the nova ejecta even when accreted material has (e) a solar abundance distribution.** |
| **(89) Models produce (eb) maximum temperature evolution profiles and nucleosynthesis yields (eb) in good agreement with (e&eb) models that generate (eb) enriched ejecta by assuming (eb) that the accreted material was pre-mixed.** |
| **(90) Authors confirm for (e) ONe novae the result we reported previously, i.e.\ we found that 3He could be produced {\it in situ} in (eb) solar-composition envelopes accreted with (e&eb) slow rates (M˙<10−10M⊙/yr) by (e) cold (TWD<107 K) CO WDs, and that convection was triggered by (eb) 3He burning before (e) the nova outburst in that case.** |
| **(91) In addition, authors now find that the interplay between the 3He production and (e&eb) destruction in the solar-composition envelope accreted with (e&eb) an intermediate rate, e.g.\ M˙=10−10M⊙/yr, by the 1.15M⊙ ONe WD with (e&eb) a relatively high initial central temperature, e.g.\ TWD=15×106 K, leads to (eb) the formation of a thick radiative buffer zone that separates (e&eb) the bottom of the convective envelope from (e) the WD surface. (Abridged) Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI: 10.1093/mnras/stu1000 Cite as: arXiv:1303.6265 [astro-ph.SR] (or arXiv:1303.6265v2 [astro-ph.SR] for this version)** |
| **Hybrid C-O-Ne White Dwarfs as Progenitors of Diverse SNe Ia Pavel Denissenkov (UVic), James Truran (U. Chicago), Falk Herwig (UVic), Sam Jones (UVic), Bill Paxton (UC Santa Barbara), Ken'ichi Nomoto (U. Tokyo), Toshio Suzuki (Nihon U.), Hiroshi Toki (Osaka U.)(Submitted on 6 Nov 2014)** |
| **(92) When carbon is ignited off-center in a CO core of a super-AGB star (eb) its burning in a convective shell tends to propagate to (e&eb) the center.** |
| **(93) Whether the C flame will actually be able to reach (eb) the center depends on (e&eb) the efficiency of (e) extra mixing beneath (e&eb) the C convective shell.** |
| **(94) Whereas thermohaline mixing is too inefficient to interfere with the C-flame propagation, (eb) convective boundary mixing can prevent (e) the C burning from reaching (eb) the center.** |
| **(95) As a result, a C-O-Ne white dwarf (WD) is (=) formed, after the star has lost (e) its envelope.** |
| **(96) Such a "hybrid" WD has (e) a small CO core surrounded by (e) a thick ONe zone.** |

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| **(97) In the proposed 1D stellar evolution computations the hybrid WD is allowed to (eb) accrete C-rich material, as if it were in a close binary system and (e&eb) accreted H-rich material from (e) its companion with (e&eb) a sufficiently high rate at which the accreted H would be processed into (e&eb) He under (e&eb) stationary conditions, assuming (eb) that He could then be transformed into (e&eb) C.** |
| **(98) When the mass of the accreting WD approaches (eb) the Chandrasekhar limit, they find a series of convective Urca shell flashes associated with (e&eb) high abundances of 23Na and 25Mg.** |
| **(99) They are followed by off-center C ignition leading to (eb) convection that occupies (e) almost the entire star.** |
| **(100) To model the Urca processes, authors use the most recent well-resolved data for (e) their reaction and (e&eb) neutrino-energy loss rates.** |
| **(101) Because of the emphasized uncertainty of the convective Urca process in hybrid WD models of (e) SN Ia progenitors, authors consider a number of their potentially possible alternative instances for (e) different mixing assumptions, all of which reach (eb) a phase of explosive C ignition, either off or in (eb) the center.** |
| **(102) Our hybrid SN Ia progenitor models have (e) much lower C to O abundance ratios at (eb) the moment of the explosive C ignition than (e) their pure CO counterparts, which may explain (eb) the observed diversity of the SNe Ia. subjects: Solar and Stellar Astrophysics (astro-ph.SR)cite as: arXiv:1411.1471 [astro-ph.SR](or arXiv:1411.1471v1 [astro-ph.SR] for this version)** |
| **Rejecting proposed dense-matter equations of state with quiescent low-mass X-ray binaries Sebastien Guillot, Robert E. Rutledge (Submitted on 15 Sep 2014 (v1), last revised 20 Oct 2014 (this version, v2))** |
| **(103) Neutrons stars are unique laboratories to (e) discriminate between the various proposed equations of (e) state of matter at and above nuclear density.** |
| **(104) One sub-class of neutron stars - those inside (eb) quiescent low-mass X-ray binaries (qLMXBs) - produce (eb) a thermal surface emission from (e) which the neutron star radius (R_NS) can be measured, using (e) the widely accepted observational scenario for qLMXBs, assuming (eb) unmagnetized H atmospheres.** |
| **(105) In a combined spectral analysis, this work first reproduces (eb) a previously published measurement of the \rns, assumed to be (=) the same for all neutron stars, using (e) a slightly expanded data set.** |
| **(106) The radius measured is R_NS = 9.4 +/-1.2 km. On the basis of spectral analysis alone, this measured value is not affected by imposing (e&eb) an assumption of causality in the core.** |
| **(107) However, the assumptions underlying this R_NS measurement would be falsified by (e) the observation of any neutron star with (e&eb) a mass >2.6 Msun, since radii <11 km would be rejected if (e) causality is assumed, which would exclude (eb) most of the R_NS parameter space obtained in this analysis.** |

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| **(108) Finally, this work directly tests (e&eb) a selection of dense matter equations of states: WFF1, AP4, MPA1, PAL1, MS0, and three versions of equations of state produced through (e&eb) chiral effective theory.** |
| **(109) Two of those, MSO and PAL1, are rejected at (eb) the 99% confidence level, accounting for (e) all quantifiable uncertainties, while (e) the other cannot be excluded at >99% certainty. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE) DOI: 10.1088/2041-8205/796/1/L3 Cite as: arXiv:1409.4306 [astro-ph.HE] (or arXiv:1409.4306v2 [astro-ph.HE] for this version** |

## SECTION SIX

## Is E=Mc^2 An Exclusively Relativistic Result

**INTRODUCTION—VARIABLES USED**

**Is E=mc^2 an exclusively relativistic result? Qing-Ping Ma**

(1) The mass-energy formula E=mc^2 is thought to be derived by (e) Einstein from special relativity.
(2) The present study shows that since the formula has also been derived from (e) classical physics by Einstein, it is not an exclusively relativistic result.
(3) The formula is implied by (e) the classical electromagnetic momentum P=E/c and the Newtonian definition of momentum P=mv.
(4) Like momentum P=mv, E=mc^2 applies to (e&eb) both classical physics and special relativity, if (e) relativistic mass is used in the equation.
(5) The derivation by Einstein in 1905 is logically flawed as (e) a relativistic proof and the truly relativistic formula should be E= (E0)/(1-v^2/c^2)^(1/2) derived by Laue in 1911 and Klein in 1918.
(6) If the energy measured in one reference frame is E0, it is E= (E0)/(1-v^2/c^2)^(1/2) in a reference frame moving at velocity v relative to (e&eb) the first frame. Subjects: General Physics (physics.gen-ph); History and Philosophy of Physics (physics.hist-ph) Cite as: arXiv:1509.07517 [physics.gen-ph] (or arXiv:1509.07517v1 [physics.gen-ph] for this version)

**NOTATION**

**Module One**

By demonstrating a measurement of photon recoil heating, Novotny's group has unambiguously observed (e) the first signatures of quantum phenomena on (e&eb) optically levitated nanoparticles

$G_{13}$ : Category one of first signatures of quantum phenomena on (e&eb) optically levitated nanoparticles

$G_{14}$ : Category two of SAS

$G_{15}$ : Category three of SAS

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| $T_{13}$ : Category one of measurement of photon recoil heating, Novotny's group has unambiguously observed<br>$T_{14}$ : Category two of SAS<br>$T_{15}$ : Category three of SAS | |
| **Module Two**<br>By demonstrating a measurement of photon recoil heating, Novotny's group has unambiguously observed the first signatures of quantum phenomena on (e&eb) optically levitated nanoparticles | |
| $G_{16}$ : Category one of **demonstrating a measurement of photon recoil heating, Novotny's group has unambiguously observed the first signatures of quantum phenomena;** optically levitated nanoparticles<br>$G_{17}$ : Category two of SAS<br>$G_{18}$ : Category three of SAS | |
| $T_{16}$ : Category one of optically levitated nanoparticles **;demonstrating a measurement of photon recoil heating, Novotny's group has unambiguously observed the first signatures of quantum phenomena**<br>$T_{17}$ : Category two of SAS<br>$T_{18}$ : Category three of SAS | |
| **Module three**<br>The mass-energy formula E=mc^2 is thought to be derived by (e) Einstein from special relativity | |
| $G_{20}$ : Category one of special relativity<br>$G_{21}$ : Category two of SAS<br>$G_{22}$ : Category three of SAS | |
| $T_{20}$ : Category one of mass-energy formula E=mc^2 is thought to be derived<br>$T_{21}$ : Category two of SAS<br>$T_{22}$ : Category three of SAS | |
| **Module four**<br>The present study shows that since the formula has also been derived from (e) classical physics by Einstein, it is not an exclusively relativistic result | |
| $G_{24}$ : Category one of classical physics by Einstein, it is not an exclusively relativistic result<br>$G_{25}$ : Category two of SAS<br>$G_{26}$ : Category three of SAS | |

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| $T_{24}$ : Category one of formula has also been derived<br><br>$T_{25}$ : Category two of SAS<br><br>$T_{26}$ : Category three of SAS | |
| **Module five** | |
| $G_{28}$ : Category one of relativistic result<br><br>$G_{29}$ : Category two of SAS<br><br>$G_{30}$ : Category three of SAS | |
| $T_{28}$ : Category one of E=mc^2<br><br>$T_{29}$ : Category two of SAS<br><br>$T_{30}$ : Category three of SAS | |
| **Module six**<br><br>The formula is implied by (e) the classical electromagnetic momentum P=E/c and the Newtonian definition of momentum P=mv | |
| $G_{32}$ : Category one of classical electromagnetic momentum P=E/c and the Newtonian definition of momentum P=mv<br><br>$G_{33}$ : Category two of SAS<br><br>$G_{34}$ : Category three of SAS | |
| $T_{32}$ : Category one of formula is implied<br><br>$T_{33}$ : Category two of SAS<br><br>$T_{34}$ : Category three of SAS | |
| **Module seven**<br><br>The formula is implied by the classical electromagnetic momentum P=E/c and the Newtonian definition of momentum P=mv | |
| $G_{36}$ : Category one of **formula is implied by the classical electromagnetic momentum P=E/c;** Newtonian definition of momentum P=mv<br><br>$G_{37}$ : Category two of SAS<br><br>$G_{38}$ : Category three of SAS | |

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| $T_{36}$ : Category one of Newtonian definition of momentum P=mv **;formula is implied by the classical electromagnetic momentum P=E/c**<br><br>$T_{37}$ : Category two of SAS<br><br>$T_{38}$ : Category three of SAS | |
| **Module eight**<br><br>Like momentum P=mv, E=mc^2 applies to (e&eb) both classical physics and special relativity, if (e) relativistic mass is used in the equation | |
| $G_{40}$ : Category one of **momentum P=mv, E=mc^2 applies;** classical physics and special relativity, if (e) relativistic mass is used in the equation<br><br>$G_{41}$ : Category two of SAS<br><br>$G_{42}$ : Category three of SAS | |
| $T_{40}$ : Category one of classical physics and special relativity, if (e) relativistic mass is used in the equation **;momentum P=mv, E=mc^2 applies**<br><br>$T_{41}$ : Category two of SAS<br><br>$T_{42}$ : Category three of SAS | |
| **Module Nine**<br><br>Like momentum P=mv, E=mc^2 applies to both classical physics and special relativity, if (e) relativistic mass is used in the equation | |
| $G_{44}$ : Category one of relativistic mass is used in the equation<br><br>$G_{45}$ : Category two of SAS<br><br>$G_{46}$ : Category three of SAS | |
| $T_{44}$ : Category one of momentum P=mv, E=mc^2 applies to both classical physics and special relativity<br><br>$T_{45}$ : Category two of SAS<br><br>$T_{46}$ : Category three of SAS | |
| **The Coefficients:** | |
| $(a_{13})^{(1)}, (a_{14})^{(1)}, (a_{15})^{(1)}, (b_{13})^{(1)}, (b_{14})^{(1)}, (b_{15})^{(1)}$ $(a_{16})^{(2)}, (a_{17})^{(2)}, (a_{18})^{(2)}$ $(b_{16})^{(2)}, (b_{17})^{(2)}, (b_{18})^{(2)}$:<br>$(a_{20})^{(3)}, (a_{21})^{(3)}, (a_{22})^{(3)}$ , $(b_{20})^{(3)}, (b_{21})^{(3)}, (b_{22})^{(3)}$<br>$(a_{24})^{(4)}, (a_{25})^{(4)}, (a_{26})^{(4)}, (b_{24})^{(4)}, (b_{25})^{(4)}, (b_{26})^{(4)}, (b_{28})^{(5)}, (b_{29})^{(5)}, (b_{30})^{(5)}, (a_{28})^{(5)}, (a_{29})^{(5)}, (a_{30})^{(5)},$<br>$(a_{32})^{(6)}, (a_{33})^{(6)}, (a_{34})^{(6)}, (b_{32})^{(6)}, (b_{33})^{(6)}, (b_{34})^{(6)}$<br>$(a_{36})^{(7)}, (a_{37})^{(7)}, (a_{38})^{(7)}, (b_{36})^{(7)}, (b_{37})^{(7)}, (b_{38})^{(7)}$<br>$(a_{40})^{(8)}, (a_{41})^{(8)}, (a_{42})^{(8)}, (b_{40})^{(8)}, (b_{41})^{(8)}, (b_{42})^{(8)}$ | |

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| $(a_{44})^{(9)}, (a_{45})^{(9)}, (a_{46})^{(9)}, (b_{44})^{(9)}, (b_{45})^{(9)}, (b_{46})^{(9)}$<br><br>are Accentuation coefficients<br>$(a_{13}^{'})^{(1)}, (a_{14}^{'})^{(1)}, (a_{15}^{'})^{(1)}, (b_{13}^{'})^{(1)}, (b_{14}^{'})^{(1)}, (b_{15}^{'})^{(1)}, (a_{16}^{'})^{(2)}, (a_{17}^{'})^{(2)}, (a_{18}^{'})^{(2)}, \ (b_{16}^{'})^{(2)}, (b_{17}^{'})^{(2)}, (b_{18}^{'})^{(2)}$<br>$, (a_{20}^{'})^{(3)}, (a_{21}^{'})^{(3)}, (a_{22}^{'})^{(3)}, (b_{20}^{'})^{(3)}, (b_{21}^{'})^{(3)}, (b_{22}^{'})^{(3)}$<br>$(a_{24}^{'})^{(4)}, (a_{25}^{'})^{(4)}, (a_{26}^{'})^{(4)}, (b_{24}^{'})^{(4)}, (b_{25}^{'})^{(4)}, (b_{26}^{'})^{(4)}, (b_{28}^{'})^{(5)}, (b_{29}^{'})^{(5)}, (b_{30}^{'})^{(5)} \ (a_{28}^{'})^{(5)}, (a_{29}^{'})^{(5)}, (a_{30}^{'})^{(5)}$ ,<br>$(a_{32}^{'})^{(6)}, (a_{33}^{'})^{(6)}, (a_{34}^{'})^{(6)}, (b_{32}^{'})^{(6)}, (b_{33}^{'})^{(6)}, (b_{34}^{'})^{(6)}$<br>$(a_{36}^{'})^{(7)}, (a_{37}^{'})^{(7)}, (a_{38}^{'})^{(7)}, (b_{36}^{'})^{(7)}, (b_{37}^{'})^{(7)}, (b_{38}^{'})^{(7)}$,<br>$(a_{40}^{'})^{(8)}, (a_{41}^{'})^{(8)}, (a_{42}^{'})^{(8)}, (b_{40}^{'})^{(8)}, (b_{41}^{'})^{(8)}, (b_{42}^{'})^{(8)}$,<br>$(a_{44}^{'})^{(9)}, (a_{45}^{'})^{(9)}, (a_{46}^{'})^{(9)}, (b_{44}^{'})^{(9)}, (b_{45}^{'})^{(9)}, (b_{46}^{'})^{(9)}$,<br>are Dissipation coefficients | |
| **Module Numbered One**<br><br>**The differential system of this model is now (Module Numbered one)** | |
| $\frac{dG_{13}}{dt} = (a_{13})^{(1)}G_{14} - \left[(a_{13}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14}, t)\right]G_{13}$ | 1 |
| $\frac{dG_{14}}{dt} = (a_{14})^{(1)}G_{13} - \left[(a_{14}^{'})^{(1)} + (a_{14}^{''})^{(1)}(T_{14}, t)\right]G_{14}$ | 2 |
| $\frac{dG_{15}}{dt} = (a_{15})^{(1)}G_{14} - \left[(a_{15}^{'})^{(1)} + (a_{15}^{''})^{(1)}(T_{14}, t)\right]G_{15}$ | 3 |
| $\frac{dT_{13}}{dt} = (b_{13})^{(1)}T_{14} - \left[(b_{13}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G, t)\right]T_{13}$ | 4 |
| $\frac{dT_{14}}{dt} = (b_{14})^{(1)}T_{13} - \left[(b_{14}^{'})^{(1)} - (b_{14}^{''})^{(1)}(G, t)\right]T_{14}$ | 5 |
| $\frac{dT_{15}}{dt} = (b_{15})^{(1)}T_{14} - \left[(b_{15}^{'})^{(1)} - (b_{15}^{''})^{(1)}(G, t)\right]T_{15}$ | 6 |
| $+(a_{13}^{''})^{(1)}(T_{14}, t) =$ First augmentation factor | |
| $-(b_{13}^{''})^{(1)}(G, t) =$ First detritions factor | |
| **Module Numbered Two**<br><br>**The differential system of this model is now ( Module numbered two)** | |
| $\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \left[(a_{16}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17}, t)\right]G_{16}$ | 7 |
| $\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \left[(a_{17}^{'})^{(2)} + (a_{17}^{''})^{(2)}(T_{17}, t)\right]G_{17}$ | 8 |
| $\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \left[(a_{18}^{'})^{(2)} + (a_{18}^{''})^{(2)}(T_{17}, t)\right]G_{18}$ | 9 |
| $\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \left[(b_{16}^{'})^{(2)} - (b_{16}^{''})^{(2)}\left((G_{19}), t\right)\right]T_{16}$ | 10 |
| $\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \left[(b_{17}^{'})^{(2)} - (b_{17}^{''})^{(2)}\left((G_{19}), t\right)\right]T_{17}$ | 11 |
| $\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \left[(b_{18}^{'})^{(2)} - (b_{18}^{''})^{(2)}\left((G_{19}), t\right)\right]T_{18}$ | 12 |
| $+(a_{16}^{''})^{(2)}(T_{17}, t) =$ First augmentation factor | |
| $-(b_{16}^{''})^{(2)}\left((G_{19}), t\right) =$ First detritions factor | |
| **Module Numbered Three**<br><br>**The differential system of this model is now (Module numbered three)** | |
| $\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \left[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21}, t)\right]G_{20}$ | 13 |
| $\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \left[(a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)}(T_{21}, t)\right]G_{21}$ | 14 |

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| $\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \left[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21},t)\right]G_{22}$ | 15 |
| $\frac{dT_{20}}{dt} = (b_{20})^{(3)}T_{21} - \left[(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23},t)\right]T_{20}$ | 16 |
| $\frac{dT_{21}}{dt} = (b_{21})^{(3)}T_{20} - \left[(b_{21}^{'})^{(3)} - (b_{21}^{''})^{(3)}(G_{23},t)\right]T_{21}$ | 17 |
| $\frac{dT_{22}}{dt} = (b_{22})^{(3)}T_{21} - \left[(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23},t)\right]T_{22}$ | 18 |
| $+(a_{20}^{''})^{(3)}(T_{21},t) =$ First augmentation factor | |
| $-(b_{20}^{''})^{(3)}(G_{23},t) =$ First detritions factor | |
| **Module Numbered Four**<br><br>**The differential system of this model is now (Module numbered Four)** | |
| $\frac{dG_{24}}{dt} = (a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25},t)\right]G_{24}$ | 19 |
| $\frac{dG_{25}}{dt} = (a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25},t)\right]G_{25}$ | 20 |
| $\frac{dG_{26}}{dt} = (a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25},t)\right]G_{26}$ | 21 |
| $\frac{dT_{24}}{dt} = (b_{24})^{(4)}T_{25} - \left[(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}\left((G_{27}),t\right)\right]T_{24}$ | 22 |
| $\frac{dT_{25}}{dt} = (b_{25})^{(4)}T_{24} - \left[(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}\left((G_{27}),t\right)\right]T_{25}$ | 23 |
| $\frac{dT_{26}}{dt} = (b_{26})^{(4)}T_{25} - \left[(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}\left((G_{27}),t\right)\right]T_{26}$ | 24 |
| $+(a_{24}^{''})^{(4)}(T_{25},t) =$ **First augmentation factor** | |
| $-(b_{24}^{''})^{(4)}\left((G_{27}),t\right) =$ **First detritions factor** | |
| **Module Numbered Five**:<br><br>**The differential system of this model is now (Module number five)** | |
| $\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29},t)\right]G_{28}$ | 25 |
| $\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29},t)\right]G_{29}$ | 26 |
| $\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \left[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29},t)\right]G_{30}$ | 27 |
| $\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \left[(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}\left((G_{31}),t\right)\right]T_{28}$ | 28 |
| $\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \left[(b_{29}^{'})^{(5)} - (b_{29}^{''})^{(5)}\left((G_{31}),t\right)\right]T_{29}$ | 29 |
| $\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \left[(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}\left((G_{31}),t\right)\right]T_{30}$ | 30 |
| $+(a_{28}^{''})^{(5)}(T_{29},t) =$ **First augmentation factor** | |
| $-(b_{28}^{''})^{(5)}\left((G_{31}),t\right) =$ **First detritions factor** | |
| **Module Numbered Six**<br><br>**The differential system of this model is now (Module numbered Six)** | |
| $\frac{dG_{32}}{dt} = (a_{32})^{(6)}G_{33} - \left[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33},t)\right]G_{32}$ | 31 |
| $\frac{dG_{33}}{dt} = (a_{33})^{(6)}G_{32} - \left[(a_{33}^{'})^{(6)} + (a_{33}^{''})^{(6)}(T_{33},t)\right]G_{33}$ | 32 |
| $\frac{dG_{34}}{dt} = (a_{34})^{(6)}G_{33} - \left[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33},t)\right]G_{34}$ | 33 |
| $\frac{dT_{32}}{dt} = (b_{32})^{(6)}T_{33} - \left[(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}\left((G_{35}),t\right)\right]T_{32}$ | 34 |
| $\frac{dT_{33}}{dt} = (b_{33})^{(6)}T_{32} - \left[(b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}\left((G_{35}),t\right)\right]T_{33}$ | 35 |

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| $\frac{dT_{34}}{dt} = (b_{34})^{(6)}T_{33} - \left[(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}((G_{35}),t)\right]T_{34}$ | 36 |
| $+(a_{32}^{''})^{(6)}(T_{33},t) =$ **First augmentation factor** | |
| **Module Numbered Seven:**<br><br>**The differential system of this model is now (Seventh Module)** | |
| $\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \left[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37},t)\right]G_{36}$ | 37 |
| $\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \left[(a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37},t)\right]G_{37}$ | *38* |
| $\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \left[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37},t)\right]G_{38}$ | *39* |
| $\frac{dT_{36}}{dt} = (b_{36})^{(7)}T_{37} - \left[(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}((G_{39}),t)\right]T_{36}$ | *40* |
| $\frac{dT_{37}}{dt} = (b_{37})^{(7)}T_{36} - \left[(b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}((G_{39}),t)\right]T_{37}$ | *41* |
| $\frac{dT_{38}}{dt} = (b_{38})^{(7)}T_{37} - \left[(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}((G_{39}),t)\right]T_{38}$ | *42* |
| $+(a_{36}^{''})^{(7)}(T_{37},t) =$ **First augmentation factor** | |
| **Module Numbered Eight**<br><br>The differential system of this model is now | |
| $\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \left[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41},t)\right]G_{40}$ | *43* |
| $\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \left[(a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41},t)\right]G_{41}$ | *44* |
| $\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \left[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41},t)\right]G_{42}$ | *45* |
| $\frac{dT_{40}}{dt} = (b_{40})^{(8)}T_{41} - \left[(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}((G_{43}),t)\right]T_{40}$ | *46* |
| $\frac{dT_{41}}{dt} = (b_{41})^{(8)}T_{40} - \left[(b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}((G_{43}),t)\right]T_{41}$ | *47* |
| $\frac{dT_{42}}{dt} = (b_{42})^{(8)}T_{41} - \left[(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}((G_{43}),t)\right]T_{42}$ | *48* |
| **Module Numbered Nine**<br><br>The differential system of this model is now | |
| $\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \left[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45},t)\right]G_{44}$ | 49 |
| $\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \left[(a_{45}^{'})^{(9)} + (a_{45}^{''})^{(9)}(T_{45},t)\right]G_{45}$ | 50 |
| $\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \left[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45},t)\right]G_{46}$ | 51 |
| $\frac{dT_{44}}{dt} = (b_{44})^{(9)}T_{45} - \left[(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}((G_{47}),t)\right]T_{44}$ | 52 |
| $\frac{dT_{45}}{dt} = (b_{45})^{(9)}T_{44} - \left[(b_{45}^{'})^{(9)} - (b_{45}^{''})^{(9)}((G_{47}),t)\right]T_{45}$ | 53 |
| $\frac{dT_{46}}{dt} = (b_{46})^{(9)}T_{45} - \left[(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}((G_{47}),t)\right]T_{46}$ | 54 |
| $+(a_{44}^{''})^{(9)}(T_{45},t) =$ **First augmentation factor** | |
| $-(b_{44}^{''})^{(9)}((G_{47}),t) =$ **First detrition factor** | |
| $\frac{dG_{13}}{dt} = (a_{13})^{(1)}G_{14} - \left[\begin{array}{ccc} (a_{13}^{'})^{(1)}\boxed{+(a_{13}^{''})^{(1)}(T_{14},t)} & \boxed{+(a_{16}^{''})^{(2,2,)}(T_{17},t)} & \boxed{+(a_{20}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{24}^{''})^{(4,4,4,4,)}(T_{25},t)} & \boxed{+(a_{28}^{''})^{(5,5,5,5,)}(T_{29},t)} & \boxed{+(a_{32}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{36}^{''})^{(7,7)}(T_{37},t)} & \boxed{+(a_{40}^{''})^{(8,8)}(T_{41},t)} & \boxed{+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{array}\right]G_{13}$ | *55* |

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| $\frac{dG_{14}}{dt} = (a_{14})^{(1)}G_{13} - \begin{bmatrix} (a_{14}^{'})^{(1)}\boxed{+(a_{14}^{''})^{(1)}(T_{14},t)}\boxed{+(a_{17}^{''})^{(2,2,)}(T_{17},t)}\boxed{+(a_{21}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{25}^{''})^{(4,4,4,4,)}(T_{25},t)}\boxed{+(a_{29}^{''})^{(5,5,5,5,)}(T_{29},t)}\boxed{+(a_{33}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{37}^{''})^{(7,7)}(T_{37},t)}\boxed{+(a_{41}^{''})^{(8,8)}(T_{41},t)}\boxed{+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{14}$ | *56* |
| $\frac{dG_{15}}{dt} = (a_{15})^{(1)}G_{14} - \begin{bmatrix} (a_{15}^{'})^{(1)}\boxed{+(a_{15}^{''})^{(1)}(T_{14},t)}\boxed{+(a_{18}^{''})^{(2,2,)}(T_{17},t)}\boxed{+(a_{22}^{''})^{(3,3,)}(T_{21},t)} \\ \boxed{+(a_{26}^{''})^{(4,4,4,4,)}(T_{25},t)}\boxed{+(a_{30}^{''})^{(5,5,5,5,)}(T_{29},t)}\boxed{+(a_{34}^{''})^{(6,6,6,6,)}(T_{33},t)} \\ \boxed{+(a_{38}^{''})^{(7,7)}(T_{37},t)}\boxed{+(a_{42}^{''})^{(8,8)}(T_{41},t)}\boxed{+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{15}$ | *57* |
| Where $\boxed{(a_{13}^{''})^{(1)}(T_{14},t)}$, $\boxed{(a_{14}^{''})^{(1)}(T_{14},t)}$, $\boxed{(a_{15}^{''})^{(1)}(T_{14},t)}$ are first augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{16}^{''})^{(2,2,)}(T_{17},t)}$, $\boxed{+(a_{17}^{''})^{(2,2,)}(T_{17},t)}$, $\boxed{+(a_{18}^{''})^{(2,2,)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{20}^{''})^{(3,3,)}(T_{21},t)}$, $\boxed{+(a_{21}^{''})^{(3,3,)}(T_{21},t)}$, $\boxed{+(a_{22}^{''})^{(3,3,)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{24}^{''})^{(4,4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{25}^{''})^{(4,4,4,4,)}(T_{25},t)}$, $\boxed{+(a_{26}^{''})^{(4,4,4,4,)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{28}^{''})^{(5,5,5,5,)}(T_{29},t)}$, $\boxed{+(a_{29}^{''})^{(5,5,5,5,)}(T_{29},t)}$, $\boxed{+(a_{30}^{''})^{(5,5,5,5,)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{32}^{''})^{(6,6,6,6,)}(T_{33},t)}$, $\boxed{+(a_{33}^{''})^{(6,6,6,6,)}(T_{33},t)}$, $\boxed{+(a_{34}^{''})^{(6,6,6,6,)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3<br>$\boxed{+(a_{38}^{''})^{(7,7)}(T_{37},t)}$ $\boxed{+(a_{37}^{''})^{(7,7)}(T_{37},t)}$ $\boxed{+(a_{36}^{''})^{(7,7)}(T_{37},t)}$ are seventh augmentation coefficient for 1,2,3<br>$\boxed{+(a_{40}^{''})^{(8,8)}(T_{41},t)}$ $\boxed{+(a_{41}^{''})^{(8,8)}(T_{41},t)}$ $\boxed{+(a_{42}^{''})^{(8,8)}(T_{41},t)}$ are eight augmentation coefficient for 1,2,3<br>$\boxed{+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)}$, $\boxed{+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)}$ are ninth augmentation coefficient for 1,2,3 | |
| $\frac{dT_{13}}{dt} = (b_{13})^{(1)}T_{14} - \begin{bmatrix} (b_{13}^{'})^{(1)}\boxed{-(b_{13}^{''})^{(1)}(G,t)}\boxed{-(b_{16}^{''})^{(2,2,)}(G_{19},t)}\boxed{-(b_{20}^{''})^{(3,3,)}(G_{23},t)} \\ \boxed{-(b_{24}^{''})^{(4,4,4,4,)}(G_{27},t)}\boxed{-(b_{28}^{''})^{(5,5,5,5,)}(G_{31},t)}\boxed{-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)} \\ \boxed{-(b_{36}^{''})^{(7,7,)}(G_{39},t)}\boxed{-(b_{40}^{''})^{(8,8)}(G_{43},t)}\boxed{-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{13}$ | 58 |
| $\frac{dT_{14}}{dt} = (b_{14})^{(1)}T_{13} - \begin{bmatrix} (b_{14}^{'})^{(1)}\boxed{-(b_{14}^{''})^{(1)}(G,t)}\boxed{-(b_{17}^{''})^{(2,2,)}(G_{19},t)}\boxed{-(b_{21}^{''})^{(3,3,)}(G_{23},t)} \\ \boxed{-(b_{25}^{''})^{(4,4,4,4,)}(G_{27},t)}\boxed{-(b_{29}^{''})^{(5,5,5,5,)}(G_{31},t)}\boxed{-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)} \\ \boxed{-(b_{37}^{''})^{(7,7,)}(G_{39},t)}\boxed{-(b_{41}^{''})^{(8,8)}(G_{43},t)}\boxed{-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{14}$ | 59 |
| $\frac{dT_{15}}{dt} = (b_{15})^{(1)}T_{14} - \begin{bmatrix} (b_{15}^{'})^{(1)}\boxed{-(b_{15}^{''})^{(1)}(G,t)}\boxed{-(b_{18}^{''})^{(2,2,)}(G_{19},t)}\boxed{-(b_{22}^{''})^{(3,3,)}(G_{23},t)} \\ \boxed{-(b_{26}^{''})^{(4,4,4,4,)}(G_{27},t)}\boxed{-(b_{30}^{''})^{(5,5,5,5,)}(G_{31},t)}\boxed{-(b_{34}^{''})^{(6,6,6,6,)}(G_{35},t)} \\ \boxed{-(b_{38}^{''})^{(7,7,)}(G_{39},t)}\boxed{-(b_{42}^{''})^{(8,8)}(G_{43},t)}\boxed{-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{15}$ | *60* |
| Where $\boxed{-(b_{13}^{''})^{(1)}(G,t)}$, $\boxed{-(b_{14}^{''})^{(1)}(G,t)}$, $\boxed{-(b_{15}^{''})^{(1)}(G,t)}$ are first detrition coefficients for category 1, 2 and 3 | |

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| $-(b_{16}^{''})^{(2,2,)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3<br>$-(b_{20}^{''})^{(3,3,)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3<br>$-(b_{24}^{''})^{(4,4,4,4,)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3<br>$-(b_{28}^{''})^{(5,5,5,5,)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3<br>$-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,)}(G_{35},t)$ are sixth detrition coefficients for category 1, 2 and 3<br>$-(b_{37}^{''})^{(7,7,)}(G_{39},t)$, $-(b_{36}^{''})^{(7,7,)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3<br>$-(b_{40}^{''})^{(8,8)}(G_{43},t)$ $-(b_{41}^{''})^{(8,8)}(G_{43},t)$ $-(b_{42}^{''})^{(8,8)}(G_{43},t)$ are eight detrition coefficients for category 1, 2 and 3<br>$-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3 | |
| | |
| $\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \left[ \begin{array}{c} (a_{16}^{'})^{(2)} +(a_{16}^{''})^{(2)}(T_{17},t) +(a_{13}^{''})^{(1,1,)}(T_{14},t) +(a_{20}^{''})^{(3,3,3)}(T_{21},t) \\ +(a_{24}^{''})^{(4,4,4,4,4)}(T_{25},t) +(a_{28}^{''})^{(5,5,5,5,5)}(T_{29},t) +(a_{32}^{''})^{(6,6,6,6,6)}(T_{33},t) \\ +(a_{36}^{''})^{(7,7,7)}(T_{37},t) +(a_{40}^{''})^{(8,8,8)}(T_{41},t) +(a_{44}^{''})^{(9,9)}(T_{45},t) \end{array} \right] G_{16}$ | *61* |
| $\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \left[ \begin{array}{c} (a_{17}^{'})^{(2)} +(a_{17}^{''})^{(2)}(T_{17},t) +(a_{14}^{''})^{(1,1,)}(T_{14},t) +(a_{21}^{''})^{(3,3,3)}(T_{21},t) \\ +(a_{25}^{''})^{(4,4,4,4,4)}(T_{25},t) +(a_{29}^{''})^{(5,5,5,5,5)}(T_{29},t) +(a_{33}^{''})^{(6,6,6,6,6)}(T_{33},t) \\ +(a_{37}^{''})^{(7,7,7)}(T_{37},t) +(a_{41}^{''})^{(8,8,8)}(T_{41},t) +(a_{45}^{''})^{(9,9)}(T_{45},t) \end{array} \right] G_{17}$ | *62* |
| $\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \left[ \begin{array}{c} (a_{18}^{'})^{(2)} +(a_{18}^{''})^{(2)}(T_{17},t) +(a_{15}^{''})^{(1,1,)}(T_{14},t) +(a_{22}^{''})^{(3,3,3)}(T_{21},t) \\ +(a_{26}^{''})^{(4,4,4,4,4)}(T_{25},t) +(a_{30}^{''})^{(5,5,5,5,5)}(T_{29},t) +(a_{34}^{''})^{(6,6,6,6,6)}(T_{33},t) \\ +(a_{38}^{''})^{(7,7,7)}(T_{37},t) +(a_{42}^{''})^{(8,8,8)}(T_{41},t) +(a_{46}^{''})^{(9,9)}(T_{45},t) \end{array} \right] G_{18}$ | *63* |
| Where $+(a_{16}^{''})^{(2)}(T_{17},t)$, $+(a_{17}^{''})^{(2)}(T_{17},t)$, $+(a_{18}^{''})^{(2)}(T_{17},t)$ are first augmentation coefficients for category 1, 2 and 3<br>$+(a_{13}^{''})^{(1,1,)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,)}(T_{14},t)$ are second augmentation coefficient for category 1, 2 and 3<br>$+(a_{20}^{''})^{(3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3<br>$+(a_{24}^{''})^{(4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3<br>$+(a_{28}^{''})^{(5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficient for category 1, 2 and 3 | |

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| $+(a_{32}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation coefficient for category 1, 2 and 3<br>$+(a_{36}^{''})^{(7,7,7)}(T_{37},t)$, $+(a_{37}^{''})^{(7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7)}(T_{37},t)$ are seventh augmentation coefficient for category 1, 2 and 3<br>$+(a_{40}^{''})^{(8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8)}(T_{41},t)$ are eight augmentation coefficient for category 1, 2 and 3<br>$+(a_{44}^{''})^{(9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9)}(T_{45},t)$, $+(a_{46}^{''})^{(9,9)}(T_{45},t)$ are ninth augmentation coefficient for category 1, 2 and 3 | |
| $\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \begin{bmatrix} (b_{16}^{'})^{(2)}(-(b_{16}^{''})^{(2)}(G_{19},t)) & (-(b_{13}^{''})^{(1,1,)}(G,t)) & (-(b_{20}^{''})^{(3,3,3,)}(G_{23},t)) \\ (-(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t)) & (-(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t)) & (-(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t)) \\ (-(b_{36}^{''})^{(7,7,7)}(G_{39},t)) & (-(b_{40}^{''})^{(8,8,8)}(G_{43},t)) & (-(b_{44}^{''})^{(9,9)}(G_{47},t)) \end{bmatrix} T_{16}$ | *64* |
| $\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \begin{bmatrix} (b_{17}^{'})^{(2)}(-(b_{17}^{''})^{(2)}(G_{19},t)) & (-(b_{14}^{''})^{(1,1,)}(G,t)) & (-(b_{21}^{''})^{(3,3,3,)}(G_{23},t)) \\ (-(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t)) & (-(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t)) & (-(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t)) \\ (-(b_{37}^{''})^{(7,7,7)}(G_{39},t)) & (-(b_{41}^{''})^{(8,8,8)}(G_{43},t)) & (-(b_{45}^{''})^{(9,9)}(G_{47},t)) \end{bmatrix} T_{17}$ | *65* |
| $\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \begin{bmatrix} (b_{18}^{'})^{(2)}(-(b_{18}^{''})^{(2)}(G_{19},t)) & (-(b_{15}^{''})^{(1,1,)}(G,t)) & (-(b_{22}^{''})^{(3,3,3,)}(G_{23},t)) \\ (-(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t)) & (-(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t)) & (-(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t)) \\ (-(b_{38}^{''})^{(7,7,7)}(G_{39},t)) & (-(b_{42}^{''})^{(8,8,8)}(G_{43},t)) & (-(b_{46}^{''})^{(9,9)}(G_{47},t)) \end{bmatrix} T_{18}$ | *66* |
| where $-(b_{16}^{''})^{(2)}(G_{19},t)$, $-(b_{17}^{''})^{(2)}(G_{19},t)$, $-(b_{18}^{''})^{(2)}(G_{19},t)$ are first detrition coefficients for category 1, 2 and 3<br>$-(b_{13}^{''})^{(1,1,)}(G,t)$, $-(b_{14}^{''})^{(1,1,)}(G,t)$, $-(b_{15}^{''})^{(1,1,)}(G,t)$ are second detrition coefficients for category 1,2 and 3<br>$-(b_{20}^{''})^{(3,3,3,)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,)}(G_{23},t)$ are third detrition coefficients for category 1,2 and 3<br>$-(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1,2 and 3<br>$-(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1,2 and 3<br>$-(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$ are sixth detrition coefficients for category 1,2 and 3<br>$-(b_{36}^{''})^{(7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1,2 and 3<br>$-(b_{40}^{''})^{(8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8)}(G_{43},t)$ are eight detrition coefficients for category 1,2 and 3<br>$-(b_{44}^{''})^{(9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1,2 and 3 | |

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| $\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \left[\begin{array}{c} (a_{20}^{'})^{(3)}\boxed{+(a_{20}^{''})^{(3)}(T_{21},t)}\ \boxed{+(a_{16}^{''})^{(2,2,2)}(T_{17},t)}\ \boxed{+(a_{13}^{''})^{(1,1,1,)}(T_{14},t)} \\ \boxed{+(a_{24}^{''})^{(4,4,4,4,4,4)}(T_{25},t)}\ \boxed{+(a_{28}^{''})^{(5,5,5,5,5,5)}(T_{29},t)}\ \boxed{+(a_{32}^{''})^{(6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{36}^{''})^{(7,7,7,7)}(T_{37},t)}\ \boxed{+(a_{40}^{''})^{(8,8,8,8)}(T_{41},t)}\ \boxed{+(a_{44}^{''})^{(9,9,9)}(T_{45},t)} \end{array}\right] G_{20}$ | 67 |
| $\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \left[\begin{array}{c} (a_{21}^{'})^{(3)}\boxed{+(a_{21}^{''})^{(3)}(T_{21},t)}\ \boxed{+(a_{17}^{''})^{(2,2,2)}(T_{17},t)}\ \boxed{+(a_{14}^{''})^{(1,1,1,)}(T_{14},t)} \\ \boxed{+(a_{25}^{''})^{(4,4,4,4,4,4)}(T_{25},t)}\ \boxed{+(a_{29}^{''})^{(5,5,5,5,5,5)}(T_{29},t)}\ \boxed{+(a_{33}^{''})^{(6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{37}^{''})^{(7,7,7,7)}(T_{37},t)}\ \boxed{+(a_{41}^{''})^{(8,8,8,8)}(T_{41},t)}\ \boxed{+(a_{45}^{''})^{(9,9,9)}(T_{45},t)} \end{array}\right] G_{21}$ | 68 |
| $\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \left[\begin{array}{c} (a_{22}^{'})^{(3)}\boxed{+(a_{22}^{''})^{(3)}(T_{21},t)}\ \boxed{+(a_{18}^{''})^{(2,2,2)}(T_{17},t)}\ \boxed{+(a_{15}^{''})^{(1,1,1,)}(T_{14},t)} \\ \boxed{+(a_{26}^{''})^{(4,4,4,4,4,4)}(T_{25},t)}\ \boxed{+(a_{30}^{''})^{(5,5,5,5,5,5)}(T_{29},t)}\ \boxed{+(a_{34}^{''})^{(6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{38}^{''})^{(7,7,7,7)}(T_{37},t)}\ \boxed{+(a_{42}^{''})^{(8,8,8,8)}(T_{41},t)}\ \boxed{+(a_{46}^{''})^{(9,9,9)}(T_{45},t)} \end{array}\right] G_{22}$ | 69 |
| $\boxed{+(a_{20}^{''})^{(3)}(T_{21},t)}$, $\boxed{+(a_{21}^{''})^{(3)}(T_{21},t)}$, $\boxed{+(a_{22}^{''})^{(3)}(T_{21},t)}$ are first augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{16}^{''})^{(2,2,2)}(T_{17},t)}$, $\boxed{+(a_{17}^{''})^{(2,2,2)}(T_{17},t)}$, $\boxed{+(a_{18}^{''})^{(2,2,2)}(T_{17},t)}$ are second augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{13}^{''})^{(1,1,1,)}(T_{14},t)}$, $\boxed{+(a_{14}^{''})^{(1,1,1,)}(T_{14},t)}$, $\boxed{+(a_{15}^{''})^{(1,1,1,)}(T_{14},t)}$ are third augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{24}^{''})^{(4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{25}^{''})^{(4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{26}^{''})^{(4,4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{28}^{''})^{(5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{29}^{''})^{(5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{30}^{''})^{(5,5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{32}^{''})^{(6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{33}^{''})^{(6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{34}^{''})^{(6,6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{36}^{''})^{(7,7,7,7)}(T_{37},t)}$, $\boxed{+(a_{37}^{''})^{(7,7,7,7)}(T_{37},t)}$, $\boxed{+(a_{38}^{''})^{(7,7,7,7)}(T_{37},t)}$ are seventh augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{40}^{''})^{(8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{41}^{''})^{(8,8,8,8)}(T_{41},t)}$, $\boxed{+(a_{42}^{''})^{(8,8,8,8)}(T_{41},t)}$ are eight augmentation coefficients for category 1, 2 and 3<br>$\boxed{+(a_{44}^{''})^{(9,9,9)}(T_{45},t)}$, $\boxed{+(a_{45}^{''})^{(9,9,9)}(T_{45},t)}$, $\boxed{+(a_{46}^{''})^{(9,9,9)}(T_{45},t)}$ are ninth augmentation coefficients for category 1, 2 and 3 | |
| $\frac{dT_{20}}{dt} = (b_{20})^{(3)}T_{21} - \left[\begin{array}{c} (b_{20}^{'})^{(3)}\boxed{-(b_{20}^{''})^{(3)}(G_{23},t)}\ \boxed{-(b_{16}^{''})^{(2,2,2)}(G_{19},t)}\ \boxed{-(b_{13}^{''})^{(1,1,1,)}(G,t)} \\ \boxed{-(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t)}\ \boxed{-(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t)}\ \boxed{-(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{36}^{''})^{(7,7,7,7)}(G_{39},t)}\ \boxed{-(b_{40}^{''})^{(8,8,8,8)}(G_{43},t)}\ \boxed{-(b_{44}^{''})^{(9,9,9)}(G_{47},t)} \end{array}\right] T_{20}$ | 70 |
| $\frac{dT_{21}}{dt} = (b_{21})^{(3)}T_{20} - \left[\begin{array}{c} (b_{21}^{'})^{(3)}\boxed{-(b_{21}^{''})^{(3)}(G_{23},t)}\ \boxed{-(b_{17}^{''})^{(2,2,2)}(G_{19},t)}\ \boxed{-(b_{14}^{''})^{(1,1,1,)}(G,t)} \\ \boxed{-(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t)}\ \boxed{-(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t)}\ \boxed{-(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{37}^{''})^{(7,7,7,7)}(G_{39},t)}\ \boxed{-(b_{41}^{''})^{(8,8,8,8)}(G_{43},t)}\ \boxed{-(b_{45}^{''})^{(9,9,9)}(G_{47},t)} \end{array}\right] T_{21}$ | 71 |

$$\frac{dT_{22}}{dt} = (b_{22})^{(3)}T_{21} - \begin{bmatrix} (b'_{22})^{(3)}\boxed{-(b''_{22})^{(3)}(G_{23},t)} & \boxed{-(b''_{18})^{(2,2,2)}(G_{19},t)} & \boxed{-(b''_{15})^{(1,1,1,)}(G,t)} \\ \boxed{-(b''_{26})^{(4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{38})^{(7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9,9)}(G_{47},t)} \end{bmatrix} T_{22} \quad (72)$$

$\boxed{-(b''_{20})^{(3)}(G_{23},t)}$, $\boxed{-(b''_{21})^{(3)}(G_{23},t)}$, $\boxed{-(b''_{22})^{(3)}(G_{23},t)}$ *are first detrition coefficients for category 1, 2 and 3*

$\boxed{-(b''_{16})^{(2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{17})^{(2,2,2)}(G_{19},t)}$, $\boxed{-(b''_{18})^{(2,2,2)}(G_{19},t)}$ *are second detrition coefficients for category 1, 2 and 3*

$\boxed{-(b''_{13})^{(1,1,1,)}(G,t)}$, $\boxed{-(b''_{14})^{(1,1,1,)}(G,t)}$, $\boxed{-(b''_{15})^{(1,1,1,)}(G,t)}$ *are third detrition coefficients for category 1,2 and 3*

$\boxed{-(b''_{24})^{(4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{25})^{(4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{26})^{(4,4,4,4,4,4)}(G_{27},t)}$ *are fourth detrition coefficients for category 1, 2 and 3*

$\boxed{-(b''_{28})^{(5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{29})^{(5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{30})^{(5,5,5,5,5,5)}(G_{31},t)}$ *are fifth detrition coefficients for category 1, 2 and 3*

$\boxed{-(b''_{32})^{(6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b''_{33})^{(6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b''_{34})^{(6,6,6,6,6,6)}(G_{35},t)}$ *are sixth detrition coefficients for category 1, 2 and 3*

$\boxed{-(b''_{36})^{(7,7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{37})^{(7,7,7,7)}(G_{39},t)}$ $\boxed{-(b''_{38})^{(7,7,7,7)}(G_{39},t)}$ are seventh detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{40})^{(8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{41})^{(8,8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{42})^{(8,8,8,8)}(G_{43},t)}$ are eight detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{46})^{(9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{45})^{(9,9,9)}(G_{47},t)}$, $\boxed{-(b''_{44})^{(9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2 and 3

$$\frac{dG_{24}}{dt} = (a_{24})^{(4)}G_{25} - \begin{bmatrix} (a'_{24})^{(4)}\boxed{+(a''_{24})^{(4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{16})^{(2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{36})^{(7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{40})^{(8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{44})^{(9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{24} \quad (73)$$

$$\frac{dG_{25}}{dt} = (a_{25})^{(4)}G_{24} - \begin{bmatrix} (a'_{25})^{(4)}\boxed{+(a''_{25})^{(4)}(T_{25},t)} & \boxed{+(a''_{29})^{(5,5,)}(T_{29},t)} & \boxed{+(a''_{33})^{(6,6)}(T_{33},t)} \\ \boxed{+(a''_{14})^{(1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{17})^{(2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{21})^{(3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{37})^{(7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{41})^{(8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{45})^{(9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{25} \quad (74)$$

$$\frac{dG_{26}}{dt} = (a_{26})^{(4)}G_{25} - \begin{bmatrix} (a'_{26})^{(4)}\boxed{+(a''_{26})^{(4)}(T_{25},t)} & \boxed{+(a''_{30})^{(5,5,)}(T_{29},t)} & \boxed{+(a''_{34})^{(6,6,)}(T_{33},t)} \\ \boxed{+(a''_{15})^{(1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{18})^{(2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{22})^{(3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{38})^{(7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{42})^{(8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{46})^{(9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{26} \quad (75)$$

$\boxed{(a''_{24})^{(4)}(T_{25},t)}$, $\boxed{(a''_{25})^{(4)}(T_{25},t)}$, $\boxed{(a''_{26})^{(4)}(T_{25},t)}$ *are first augmentation coefficients category* 1, 2 3

$\boxed{+(a''_{28})^{(5,5,)}(T_{29},t)}$, $\boxed{+(a''_{29})^{(5,5,)}(T_{29},t)}$, $\boxed{+(a''_{30})^{(5,5,)}(T_{29},t)}$ *are second augmentation coefficient for category* 1, 2 *and* 3

$\boxed{+(a''_{32})^{(6,6,)}(T_{33},t)}$, $\boxed{+(a''_{33})^{(6,6,)}(T_{33},t)}$, $\boxed{+(a''_{34})^{(6,6,)}(T_{33},t)}$ *are third augmentation coefficient for category* 1, 2 *and* 3

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| $+(a_{13}^{''})^{(1,1,1,1)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1)}(T_{14},t)$ are fourth augmentation coefficients for category 1, 2 and 3 $+(a_{16}^{''})^{(2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2)}(T_{17},t)$ are fifth augmentation coefficients for category 1, 2 and 3 $+(a_{20}^{''})^{(3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3)}(T_{21},t)$ are sixth augmentation coefficients for category 1, 2 and 3 $+(a_{36}^{''})^{(7,7,7,7,7)}(T_{37},t)$, $+(a_{37}^{''})^{(7,7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7,7)}(T_{37},t)$ are seventh augmentation coefficients for category 1, 2 and 3 $+(a_{40}^{''})^{(8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8,8)}(T_{41},t)$ are eighth augmentation coefficients for category 1, 2 and 3 $+(a_{46}^{''})^{(9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9)}(T_{45},t)$ are ninth detrition coefficients for category 1 2 3 | |
| $\frac{dT_{24}}{dt} = (b_{24})^{(4)}T_{25} - \begin{bmatrix} (b_{24}^{'})^{(4)} & -(b_{24}^{''})^{(4)}(G_{27},t) & -(b_{28}^{''})^{(5,5,)}(G_{31},t) & -(b_{32}^{''})^{(6,6,)}(G_{35},t) \\ & -(b_{13}^{''})^{(1,1,1,1)}(G,t) & -(b_{16}^{''})^{(2,2,2,2)}(G_{19},t) & -(b_{20}^{''})^{(3,3,3,3)}(G_{23},t) \\ & -(b_{36}^{''})^{(7,7,7,7,7)}(G_{39},t) & -(b_{40}^{''})^{(8,8,8,8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{24}$ | 76 |
| $\frac{dT_{25}}{dt} = (b_{25})^{(4)}T_{24} - \begin{bmatrix} (b_{25}^{'})^{(4)} & -(b_{25}^{''})^{(4)}(G_{27},t) & -(b_{29}^{''})^{(5,5,)}(G_{31},t) & -(b_{33}^{''})^{(6,6,)}(G_{35},t) \\ & -(b_{14}^{''})^{(1,1,1,1)}(G,t) & -(b_{17}^{''})^{(2,2,2,2)}(G_{19},t) & -(b_{21}^{''})^{(3,3,3,3)}(G_{23},t) \\ & -(b_{37}^{''})^{(7,7,7,7,7)}(G_{39},t) & -(b_{41}^{''})^{(8,8,8,8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{25}$ | 77 |
| $\frac{dT_{26}}{dt} = (b_{26})^{(4)}T_{25} - \begin{bmatrix} (b_{26}^{'})^{(4)} & -(b_{26}^{''})^{(4)}(G_{27},t) & -(b_{30}^{''})^{(5,5,)}(G_{31},t) & -(b_{34}^{''})^{(6,6,)}(G_{35},t) \\ & -(b_{15}^{''})^{(1,1,1,1)}(G,t) & -(b_{18}^{''})^{(2,2,2,2)}(G_{19},t) & -(b_{22}^{''})^{(3,3,3,3)}(G_{23},t) \\ & -(b_{38}^{''})^{(7,7,7,7,7)}(G_{39},t) & -(b_{42}^{''})^{(8,8,8,8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{26}$ | 78 |

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|---|---|
| $Where\ \boxed{-(b_{24}^{''})^{(4)}(G_{27},t)}, \boxed{-(b_{25}^{''})^{(4)}(G_{27},t)}, \boxed{-(b_{26}^{''})^{(4)}(G_{27},t)}\ are\ first\ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b_{28}^{''})^{(5,5,)}(G_{31},t)}, \boxed{-(b_{29}^{''})^{(5,5,)}(G_{31},t)}, \boxed{-(b_{30}^{''})^{(5,5,)}(G_{31},t)}\ are\ second\ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b_{32}^{''})^{(6,6,)}(G_{35},t)}, \boxed{-(b_{33}^{''})^{(6,6,)}(G_{35},t)}, \boxed{-(b_{34}^{''})^{(6,6,)}(G_{35},t)}\ are\ third\ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b_{13}^{''})^{(1,1,1,1)}(G,t)}, \boxed{-(b_{14}^{''})^{(1,1,1,1)}(G,t)}, \boxed{-(b_{15}^{''})^{(1,1,1,1)}(G,t)}$ $are\ fourth\ \ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b_{16}^{''})^{(2,2,2,2)}(G_{19},t)}, \boxed{-(b_{17}^{''})^{(2,2,2,2)}(G_{19},t)}, \boxed{-(b_{18}^{''})^{(2,2,2,2)}(G_{19},t)}$ $are\ fifth\ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b_{20}^{''})^{(3,3,3,3)}(G_{23},t)}, \boxed{-(b_{21}^{''})^{(3,3,3,3)}(G_{23},t)}, \boxed{-(b_{22}^{''})^{(3,3,3,3)}(G_{23},t)}$ $are\ sixth\ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b_{36}^{''})^{(7,7,7,7,7)}(G_{39},t)}, \boxed{-(b_{37}^{''})^{(7,7,7,7,7)}(G_{39},t)}, \boxed{-(b_{38}^{''})^{(7,7,7,7,7)}(G_{39},t)}$ $are\ seventh\ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b_{40}^{''})^{(8,8,8,8,8)}(G_{43},t)}, \boxed{-(b_{41}^{''})^{(8,8,8,8,8)}(G_{43},t)}, \boxed{-(b_{42}^{''})^{(8,8,8,8,8)}(G_{43},t)}$ $are\ eighth\ detrition\ coefficients\ for\ category\ 1,2\ and\ 3$ $\boxed{-(b_{46}^{''})^{(9,9,9,9)}(G_{47},t)}, \boxed{-(b_{45}^{''})^{(9,9,9,9)}(G_{47},t)}, \boxed{-(b_{44}^{''})^{(9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1 2 3 | |
| $\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \begin{bmatrix} (a_{28}^{'})^{(5)}\boxed{+(a_{28}^{''})^{(5)}(T_{29},t)} & \boxed{+(a_{24}^{''})^{(4,4,)}(T_{25},t)} & \boxed{+(a_{32}^{''})^{(6,6,6)}(T_{33},t)} \\ \boxed{+(a_{13}^{''})^{(1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{16}^{''})^{(2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{20}^{''})^{(3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{36}^{''})^{(7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{40}^{''})^{(8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{44}^{''})^{(9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{28}$ | *79* |
| $\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \begin{bmatrix} (a_{29}^{'})^{(5)}\boxed{+(a_{29}^{''})^{(5)}(T_{29},t)} & \boxed{+(a_{25}^{''})^{(4,4,)}(T_{25},t)} & \boxed{+(a_{33}^{''})^{(6,6,6)}(T_{33},t)} \\ \boxed{+(a_{14}^{''})^{(1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{17}^{''})^{(2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{21}^{''})^{(3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{37}^{''})^{(7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{41}^{''})^{(8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{45}^{''})^{(9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{29}$ | *80* |
| $\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \begin{bmatrix} (a_{30}^{'})^{(5)}\boxed{+(a_{30}^{''})^{(5)}(T_{29},t)} & \boxed{+(a_{26}^{''})^{(4,4,)}(T_{25},t)} & \boxed{+(a_{34}^{''})^{(6,6,6)}(T_{33},t)} \\ \boxed{+(a_{15}^{''})^{(1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{18}^{''})^{(2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{22}^{''})^{(3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{38}^{''})^{(7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{42}^{''})^{(8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a_{46}^{''})^{(9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{30}$ | *81* |
| $Where\ \boxed{+(a_{28}^{''})^{(5)}(T_{29},t)}, \boxed{+(a_{29}^{''})^{(5)}(T_{29},t)}, \boxed{+(a_{30}^{''})^{(5)}(T_{29},t)}\ are\ first\ augmentation\ coefficients\ for\ category\ 1,2\ and\ 3$ $And\ \boxed{+(a_{24}^{''})^{(4,4,)}(T_{25},t)}, \boxed{+(a_{25}^{''})^{(4,4,)}(T_{25},t)}, \boxed{+(a_{26}^{''})^{(4,4,)}(T_{25},t)}\ are\ second\ augmentation\ coefficient\ for\ category\ 1,2\ and\ 3$ $\boxed{+(a_{32}^{''})^{(6,6,6)}(T_{33},t)}, \boxed{+(a_{33}^{''})^{(6,6,6)}(T_{33},t)}, \boxed{+(a_{34}^{''})^{(6,6,6)}(T_{33},t)}\ are\ third\ \ augmentation\ coefficient\ for\ category\ 1,2\ and\ 3$ $\boxed{+(a_{13}^{''})^{(1,1,1,1,1)}(T_{14},t)}, \boxed{+(a_{14}^{''})^{(1,1,1,1,1)}(T_{14},t)}, \boxed{+(a_{15}^{''})^{(1,1,1,1,1)}(T_{14},t)}$ are fourth augmentation coefficients for category 1,2, and 3 $\boxed{+(a_{16}^{''})^{(2,2,2,2,2)}(T_{17},t)}, \boxed{+(a_{17}^{''})^{(2,2,2,2,2)}(T_{17},t)}, \boxed{+(a_{18}^{''})^{(2,2,2,2,2)}(T_{17},t)}$ are fifth augmentation | |

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| coefficients for category 1,2,and 3<br>$+(a_{20}^{''})^{(3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3)}(T_{21},t)$ are sixth augmentation coefficients for category 1,2, 3<br>$+(a_{36}^{''})^{(7,7,7,7,7,7)}(T_{37},t)$ , $+(a_{37}^{''})^{(7,7,7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7,7,7)}(T_{37},t)$ are seventh augmentation coefficients for category 1,2, 3<br>$+(a_{40}^{''})^{(8,8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8,8,8)}(T_{41},t)$ are eighth augmentation coefficients for category 1,2, 3<br>$+(a_{46}^{''})^{(9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficients for category 1,2, 3 | |
| $\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \left[\begin{matrix}(b_{28}^{'})^{(5)}\,-(b_{28}^{''})^{(5)}(G_{31},t) & -(b_{24}^{''})^{(4,4,)}(G_{27},t) & -(b_{32}^{''})^{(6,6,6)}(G_{35},t) \\ -(b_{13}^{''})^{(1,1,1,1,1)}(G,t) & -(b_{16}^{''})^{(2,2,2,2,2)}(G_{19},t) & -(b_{20}^{''})^{(3,3,3,3,3)}(G_{23},t) \\ -(b_{36}^{''})^{(7,7,7,7,7,7)}(G_{39},t) & -(b_{40}^{''})^{(8,8,8,8,8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9,9,9,9)}(G_{47},t)\end{matrix}\right]T_{28}$ | *82* |
| $\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \left[\begin{matrix}(b_{29}^{'})^{(5)}\,-(b_{29}^{''})^{(5)}(G_{31},t) & -(b_{25}^{''})^{(4,4,)}(G_{27},t) & -(b_{33}^{''})^{(6,6,6)}(G_{35},t) \\ -(b_{14}^{''})^{(1,1,1,1,1)}(G,t) & -(b_{17}^{''})^{(2,2,2,2,2)}(G_{19},t) & -(b_{21}^{''})^{(3,3,3,3,3)}(G_{23},t) \\ -(b_{37}^{''})^{(7,7,7,7,7,7)}(G_{39},t) & -(b_{41}^{''})^{(8,8,8,8,8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9,9,9,9)}(G_{47},t)\end{matrix}\right]T_{29}$ | *83* |
| $\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \left[\begin{matrix}(b_{30}^{'})^{(5)}\,-(b_{30}^{''})^{(5)}(G_{31},t) & -(b_{26}^{''})^{(4,4,)}(G_{27},t) & -(b_{34}^{''})^{(6,6,6)}(G_{35},t) \\ -(b_{15}^{''})^{(1,1,1,1,1,)}(G,t) & -(b_{18}^{''})^{(2,2,2,2,2)}(G_{19},t) & -(b_{22}^{''})^{(3,3,3,3,3)}(G_{23},t) \\ -(b_{38}^{''})^{(7,7,7,7,7,7)}(G_{39},t) & -(b_{42}^{''})^{(8,8,8,8,8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9,9,9)}(G_{47},t)\end{matrix}\right]T_{30}$ | *84* |
| *where* $-(b_{28}^{''})^{(5)}(G_{31},t)$ , $-(b_{29}^{''})^{(5)}(G_{31},t)$ , $-(b_{30}^{''})^{(5)}(G_{31},t)$ *are first detrition coefficients for category* 1, 2 *and* 3<br>$-(b_{24}^{''})^{(4,4,)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,)}(G_{27},t)$ *are second detrition coefficients for category* 1,2 *and* 3<br>$-(b_{32}^{''})^{(6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6)}(G_{35},t)$ *are third detrition coefficients for category* 1,2 *and* 3<br>$-(b_{13}^{''})^{(1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1)}(G,t)$ , $-(b_{15}^{''})^{(1,1,1,1,1,)}(G,t)$ are fourth detrition coefficients for category 1,2, and 3<br>$-(b_{16}^{''})^{(2,2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2,2)}(G_{19},t)$ are fifth detrition coefficients for category 1,2, and 3<br>$-(b_{20}^{''})^{(3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3)}(G_{23},t)$ are sixth detrition coefficients for category 1,2, and 3<br>$-(b_{36}^{''})^{(7,7,7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1,2, and 3<br>$-(b_{42}^{''})^{(8,8,8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8,8,8)}(G_{43},t)$, $-(b_{40}^{''})^{(8,8,8,8,8,8)}(G_{43},t)$ are eighth detrition coefficients for category 1,2, and 3<br>$-(b_{46}^{''})^{(9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1,2, and 3 | |

$$\frac{dG_{32}}{dt} = (a_{32})^{(6)} G_{33} - \begin{bmatrix} (a'_{32})^{(6)} & +(a''_{32})^{(6)}(T_{33},t) & +(a''_{28})^{(5,5,5)}(T_{29},t) & +(a''_{24})^{(4,4,4,)}(T_{25},t) \\ +(a''_{13})^{(1,1,1,1,1,1)}(T_{14},t) & & +(a''_{16})^{(2,2,2,2,2,2)}(T_{17},t) & +(a''_{20})^{(3,3,3,3,3,3)}(T_{21},t) \\ +(a''_{36})^{(7,7,7,7,7,7,7)}(T_{37},t) & & +(a''_{40})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a''_{44})^{(9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{32} \quad (85)$$

$$\frac{dG_{33}}{dt} = (a_{33})^{(6)} G_{32} - \begin{bmatrix} (a'_{33})^{(6)} & +(a''_{33})^{(6)}(T_{33},t) & +(a''_{29})^{(5,5,5)}(T_{29},t) & +(a''_{25})^{(4,4,4,)}(T_{25},t) \\ +(a''_{14})^{(1,1,1,1,1,1)}(T_{14},t) & & +(a''_{17})^{(2,2,2,2,2,2)}(T_{17},t) & +(a''_{21})^{(3,3,3,3,3,3)}(T_{21},t) \\ +(a''_{37})^{(7,7,7,7,7,7,7)}(T_{37},t) & & +(a''_{41})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a''_{45})^{(9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{33} \quad (86)$$

$$\frac{dG_{34}}{dt} = (a_{34})^{(6)} G_{33} - \begin{bmatrix} (a'_{34})^{(6)} & +(a''_{34})^{(6)}(T_{33},t) & +(a''_{30})^{(5,5,5)}(T_{29},t) & +(a''_{26})^{(4,4,4,)}(T_{25},t) \\ +(a''_{15})^{(1,1,1,1,1,1)}(T_{14},t) & & +(a''_{18})^{(2,2,2,2,2,2)}(T_{17},t) & +(a''_{22})^{(3,3,3,3,3,3)}(T_{21},t) \\ +(a''_{38})^{(7,7,7,7,7,7,7)}(T_{37},t) & & +(a''_{42})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a''_{46})^{(9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{34} \quad (87)$$

$+(a''_{32})^{(6)}(T_{33},t)$, $+(a''_{33})^{(6)}(T_{33},t)$, $+(a''_{34})^{(6)}(T_{33},t)$ *are first augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{28})^{(5,5,5)}(T_{29},t)$, $+(a''_{29})^{(5,5,5)}(T_{29},t)$, $+(a''_{30})^{(5,5,5)}(T_{29},t)$ *are second augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{24})^{(4,4,4,)}(T_{25},t)$, $+(a''_{25})^{(4,4,4,)}(T_{25},t)$, $+(a''_{26})^{(4,4,4,)}(T_{25},t)$ *are third augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{13})^{(1,1,1,1,1,1)}(T_{14},t)$, $+(a''_{14})^{(1,1,1,1,1,1)}(T_{14},t)$, $+(a''_{15})^{(1,1,1,1,1,1)}(T_{14},t)$ - are fourth augmentation coefficients

$+(a''_{16})^{(2,2,2,2,2,2)}(T_{17},t)$, $+(a''_{17})^{(2,2,2,2,2,2)}(T_{17},t)$, $+(a''_{18})^{(2,2,2,2,2,2)}(T_{17},t)$ - fifth augmentation coefficients

$+(a''_{20})^{(3,3,3,3,3,3)}(T_{21},t)$, $+(a''_{21})^{(3,3,3,3,3,3)}(T_{21},t)$, $+(a''_{22})^{(3,3,3,3,3,3)}(T_{21},t)$ sixth augmentation coefficients

$+(a''_{36})^{(7,7,7,7,7,7,7)}(T_{37},t)$, $+(a''_{37})^{(7,7,7,7,7,7,7)}(T_{37},t)$, $+(a''_{38})^{(7,7,7,7,7,7,7)}(T_{37},t)$ seventh augmentation coefficients

$+(a''_{40})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a''_{41})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a''_{42})^{(8,8,8,8,8,8,8)}(T_{41},t)$ Eighth augmentation coefficients

$+(a''_{44})^{(9,9,9,9,9,9)}(T_{45},t)$, $+(a''_{45})^{(9,9,9,9,9,9)}(T_{45},t)$, $+(a''_{46})^{(9,9,9,9,9,9)}(T_{45},t)$ ninth augmentation coefficients

$$\frac{dT_{32}}{dt} = (b_{32})^{(6)} T_{33} - \begin{bmatrix} (b'_{32})^{(6)} & -(b''_{32})^{(6)}(G_{35},t) & -(b''_{28})^{(5,5,5)}(G_{31},t) & -(b''_{24})^{(4,4,4,)}(G_{27},t) \\ -(b''_{13})^{(1,1,1,1,1,1)}(G,t) & & -(b''_{16})^{(2,2,2,2,2,2)}(G_{19},t) & -(b''_{20})^{(3,3,3,3,3,3)}(G_{23},t) \\ -(b''_{36})^{(7,7,7,7,7,7,7)}(G_{39},t) & & -(b''_{40})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b''_{44})^{(9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{32} \quad (88)$$

$$\frac{dT_{33}}{dt} = (b_{33})^{(6)} T_{32} - \begin{bmatrix} (b'_{33})^{(6)} & -(b''_{33})^{(6)}(G_{35},t) & -(b''_{29})^{(5,5,5)}(G_{31},t) & -(b''_{25})^{(4,4,4,)}(G_{27},t) \\ -(b''_{14})^{(1,1,1,1,1,1)}(G,t) & & -(b''_{17})^{(2,2,2,2,2,2)}(G_{19},t) & -(b''_{21})^{(3,3,3,3,3,3)}(G_{23},t) \\ -(b''_{37})^{(7,7,7,7,7,7,7)}(G_{39},t) & & -(b''_{41})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b''_{45})^{(9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{33} \quad (89)$$

$$\frac{dT_{34}}{dt} = (b_{34})^{(6)}T_{33} - \begin{bmatrix} (b'_{34})^{(6)}\boxed{-(b''_{34})^{(6)}(G_{35},t)} & \boxed{-(b''_{30})^{(5,5,5)}(G_{31},t)} & \boxed{-(b''_{26})^{(4,4,4,)}(G_{27},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{18})^{(2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{38})^{(7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{34} \quad (90)$$

$\boxed{-(b''_{32})^{(6)}(G_{35},t)}, \boxed{-(b''_{33})^{(6)}(G_{35},t)}, \boxed{-(b''_{34})^{(6)}(G_{35},t)}$ *are first detrition coefficients for category* 1, 2 *and* 3

$\boxed{-(b''_{28})^{(5,5,5)}(G_{31},t)}, \boxed{-(b''_{29})^{(5,5,5)}(G_{31},t)}, \boxed{-(b''_{30})^{(5,5,5)}(G_{31},t)}$ *are second detrition coefficients for category* 1, 2 *and* 3

$\boxed{-(b''_{24})^{(4,4,4,)}(G_{27},t)}, \boxed{-(b''_{25})^{(4,4,4,)}(G_{27},t)}, \boxed{-(b''_{26})^{(4,4,4,)}(G_{27},t)}$ *are third detrition coefficients for category* 1,2 *and* 3

$\boxed{-(b''_{13})^{(1,1,1,1,1,1)}(G,t)}, \boxed{-(b''_{14})^{(1,1,1,1,1,1)}(G,t)}, \boxed{-(b''_{15})^{(1,1,1,1,1,1)}(G,t)}$ are fourth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{16})^{(2,2,2,2,2,2)}(G_{19},t)}, \boxed{-(b''_{17})^{(2,2,2,2,2,2)}(G_{19},t)}, \boxed{-(b''_{18})^{(2,2,2,2,2,2)}(G_{19},t)}$ are fifth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{20})^{(3,3,3,3,3,3)}(G_{23},t)}, \boxed{-(b''_{21})^{(3,3,3,3,3,3)}(G_{23},t)}, \boxed{-(b''_{22})^{(3,3,3,3,3,3)}(G_{23},t)}$ are sixth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{36})^{(7,7,7,7,7,7,7)}(G_{39},t)}, \boxed{-(b''_{37})^{(7,7,7,7,7,7,7)}(G_{39},t)}, \boxed{-(b''_{38})^{(7,7,7,7,7,7,7)}(G_{39},t)}$ are seventh detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{40})^{(8,8,8,8,8,8,8,8)}(G_{43},t)}, \boxed{-(b''_{41})^{(8,8,8,8,8,8,8,8)}(G_{43},t)}, \boxed{-(b''_{42})^{(8,8,8,8,8,8,8,8)}(G_{43},t)}$ are eighth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{46})^{(9,9,9,9,9,9)}(G_{47},t)}, \boxed{-(b''_{45})^{(9,9,9,9,9,9)}(G_{47},t)}, \boxed{-(b''_{44})^{(9,9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2, and 3

$$\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \begin{bmatrix} (a'_{36})^{(7)}\boxed{+(a''_{36})^{(7)}(T_{37},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{24})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8,)}(T_{41},t)} & \boxed{+(a''_{44})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{13} \quad (91)$$

$$\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \begin{bmatrix} (a'_{37})^{(7)}\boxed{+(a''_{37})^{(7)}(T_{37},t)} & \boxed{+(a''_{17})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{21})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{25})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{29})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{33})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{41})^{(8,8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{45})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{14} \quad (92)$$

$$\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \begin{bmatrix} (a'_{38})^{(7)}\boxed{+(a''_{38})^{(7)}(T_{37},t)} & \boxed{+(a''_{18})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{22})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{26})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{30})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{34})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{15})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{42})^{(8,8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{46})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{15} \quad (93)$$

Where $(a_{36}^{''})^{(7)}(T_{37},t)$, $(a_{37}^{''})^{(7)}(T_{37},t)$, $(a_{38}^{''})^{(7)}(T_{37},t)$ are first augmentation coefficients for category 1, 2 and 3

$+(a_{16}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t)$ are second augmentation coefficient for category 1, 2 and 3

$+(a_{20}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3

$+(a_{24}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3

$+(a_{28}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficient for category 1, 2 and 3

$+(a_{32}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation coefficient for category 1, 2 and 3

$+(a_{13}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{13}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t)$ are seventh augmentation coefficient for category 1, 2 and 3

$+(a_{42}^{''})^{(8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{40}^{''})^{(8,8,8,8,8,8,8,8,)}(T_{41},t)$ are eighth augmentation coefficient for 1,2,3

$+(a_{46}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{36}}{dt} = (b_{36})^{(7)}T_{37} - \begin{bmatrix} (b_{36}^{'})^{(7)} & -(b_{36}^{''})^{(7)}(G_{39},t) & -(b_{16}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t) & -(b_{20}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t) \\ & -(b_{24}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t) & -(b_{28}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t) & -(b_{32}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t) \\ & -(b_{13}^{''})^{(1,1,1,1,1,1,1)}(G,t) & -(b_{40}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{13} \quad (94)$$

$$\frac{dT_{37}}{dt} = (b_{37})^{(7)}T_{36} - \begin{bmatrix} (b_{37}^{'})^{(7)} & -(b_{37}^{''})^{(7)}(G_{39},t) & -(b_{17}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t) & -(b_{21}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t) \\ & -(b_{25}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t) & -(b_{29}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t) & -(b_{33}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t) \\ & -(b_{14}^{''})^{(1,1,1,1,1,1,1)}(G,t) & -(b_{41}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{14}$$

$$\frac{dT_{38}}{dt} = (b_{38})^{(7)}T_{37} - \begin{bmatrix} (b_{38}^{'})^{(7)} & -(b_{38}^{''})^{(7)}(G_{39},t) & -(b_{18}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t) & -(b_{22}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t) \\ & -(b_{26}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t) & -(b_{30}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t) & -(b_{34}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t) \\ & -(b_{15}^{''})^{(1,1,1,1,1,1,1)}(G,t) & -(b_{42}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{15}$$

Where $-(b_{36}^{''})^{(7)}(G_{39},t)$, $-(b_{37}^{''})^{(7)}(G_{39},t)$, $-(b_{38}^{''})^{(7)}(G_{39},t)$ are first detrition coefficients for category 1, 2 and 3

$-(b_{16}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3

| | |
|---|---|
| $-(b_{20}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3<br>$-(b_{24}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3<br>$-(b_{28}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3<br>$-(b_{32}^{''})^{(6,6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6,6,6)}(G_{35},t)$ are sixth detrition coefficients for category 1, 2 and 3<br>$-(b_{15}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{13}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$ are seventh detrition coefficients for category 1, 2 and 3<br>$-(b_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)$ are eighth detrition coefficients for category 1, 2 and 3<br>$-(b_{46}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3 | |
| $\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \begin{bmatrix} (a_{40}^{'})^{(8)}+(a_{40}^{''})^{(8)}(T_{41},t) & +(a_{16}^{''})^{(2,2,2,2,2,2,2,2)}(T_{17},t) & +(a_{20}^{''})^{(3,3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{24}^{''})^{(4,4,4,4,4,4,4,4)}(T_{25},t) & +(a_{28}^{''})^{(5,5,5,5,5,5,5,5)}(T_{29},t) & +(a_{32}^{''})^{(6,6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{13}^{''})^{(1,1,1,1,1,1,1,1)}(T_{14},t) & +(a_{36}^{''})^{(7,7,7,7,7,7,7,7)}(T_{37},t) & +(a_{44}^{''})^{(9,9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{13}$ | *95* |
| $\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \begin{bmatrix} (a_{41}^{'})^{(8)}+(a_{41}^{''})^{(8)}(T_{41},t) & +(a_{17}^{''})^{(2,2,2,2,2,2,2,2)}(T_{17},t) & +(a_{21}^{''})^{(3,3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{25}^{''})^{(4,4,4,4,4,4,4,4)}(T_{25},t) & +(a_{29}^{''})^{(5,5,5,5,5,5,5,5)}(T_{29},t) & +(a_{33}^{''})^{(6,6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{13}^{''})^{(1,1,1,1,1,1,1,1)}(T_{14},t) & +(a_{37}^{''})^{(7,7,7,7,7,7,7,7)}(T_{37},t) & +(a_{45}^{''})^{(9,9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{14}$ | |
| $\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \begin{bmatrix} (a_{42}^{'})^{(8)}+(a_{42}^{''})^{(8)}(T_{41},t) & +(a_{18}^{''})^{(2,2,2,2,2,2,2,2)}(T_{17},t) & +(a_{22}^{''})^{(3,3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{26}^{''})^{(4,4,4,4,4,4,4,4)}(T_{25},t) & +(a_{30}^{''})^{(5,5,5,5,5,5,5,5)}(T_{29},t) & +(a_{34}^{''})^{(6,6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{15}^{''})^{(1,1,1,1,1,1,1,1)}(T_{14},t) & +(a_{38}^{''})^{(7,7,7,7,7,7,7,7)}(T_{37},t) & +(a_{46}^{''})^{(9,9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{15}$ | |
| Where $+(a_{40}^{''})^{(8)}(T_{41},t)$, $+(a_{41}^{''})^{(8)}(T_{41},t)$, $+(a_{42}^{''})^{(8)}(T_{41},t)$ are first augmentation coefficients for category 1, 2 and 3<br>$+(a_{16}^{''})^{(2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2,2,2,2)}(T_{17},t)$ are second augmentation coefficient for category 1, 2 and 3<br>$+(a_{20}^{''})^{(3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3,3,3,3)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3 | |

$+(a_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$ , $+(a_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$ , $+(a_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3

$+(a_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$ , $+(a_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficient for category 1, 2 and 3

$+(a_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$ , $+(a_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation coefficient for category 1, 2 and 3

$+(a_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)$ $+(a_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)$ $+(a_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)$ are seventh augmentation coefficient for 1,2,3

$+(a_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)$ ,$+(a_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)$ ,$+(a_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)$ are eighth augmentation coefficient for 1,2,3

$+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ ,$+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ ,$+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{40}}{dt} =$$

$$(b_{40})^{(8)}T_{41} - \begin{bmatrix} (b_{40}^{'})^{(8)}\left[-(b_{40}^{''})^{(8)}(G_{43},t)\right] & \left[-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)\right] & \left[-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)\right] \\ \left[-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)\right] & \left[-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)\right] & \left[-(b_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)\right] \\ \left[-(b_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)\right] & \left[-(b_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)\right] & \left[-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{13}$$

$$\frac{dT_{41}}{dt} =$$

$$(b_{41})^{(8)}T_{40} - \begin{bmatrix} (b_{41}^{'})^{(8)}\left[-(b_{41}^{''})^{(8)}(G_{43},t)\right] & \left[-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)\right] & \left[-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)\right] \\ \left[-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)\right] & \left[-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)\right] & \left[-(b_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)\right] \\ \left[-(b_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)\right] & \left[-(b_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)\right] & \left[-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{14}$$

$$\frac{dT_{42}}{dt} =$$

$$(b_{42})^{(8)}T_{41} - \begin{bmatrix} (b_{42}^{'})^{(8)}\left[-(b_{42}^{''})^{(8)}(G_{43},t)\right] & \left[-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)\right] & \left[-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)\right] \\ \left[-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)\right] & \left[-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)\right] & \left[-(b_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)\right] \\ \left[-(b_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)\right] & \left[-(b_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)\right] & \left[-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{15}$$

Where $-(b_{36}^{''})^{(7)}(G_{39},t)$ , $-(b_{37}^{''})^{(7)}(G_{39},t)$ , $-(b_{38}^{''})^{(7)}(G_{39},t)$ are first detrition coefficients for category 1, 2 and 3

$-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)$ , $-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)$ , $-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3

$-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)$ , $-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)$ , $-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3

$-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)$ , $-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)$ , $-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3

$-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)$ , $-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)$ , $-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3

$-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{15}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$ are sixth detrition coefficients for category 1, 2 and 3

$-(b_{13}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{38}^{''})^{(7,7,)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3

$-(b_{36}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)$ are eighth detrition coefficients for category 1, 2 and 3

$-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3

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$$\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \begin{bmatrix} (a_{44}^{'})^{(9)}\boxed{+(a_{44}^{''})^{(9)}(T_{45},t)} & \boxed{+(a_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)} \end{bmatrix} G_{13}$$

$$\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \begin{bmatrix} (a_{45}^{'})^{(9)}\boxed{+(a_{45}^{''})^{(9)}(T_{45},t)} & \boxed{+(a_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)} \end{bmatrix} G_{14}$$

$$\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \begin{bmatrix} (a_{46}^{'})^{(9)}\boxed{+(a_{46}^{''})^{(9)}(T_{37},t)} & \boxed{+(a_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)} \end{bmatrix} G_{15}$$

Where $+(a_{44}^{''})^{(9)}(T_{45},t)$, $+(a_{45}^{''})^{(9)}(T_{45},t)$, $+(a_{46}^{''})^{(9)}(T_{37},t)$ are first augmentation coefficients for category 1, 2 and 3

$+(a_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)$ are second augmentation coefficient for category 1, 2 and 3

$+(a_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3

$+(a_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3

$+(a_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficient for category 1, 2 and 3

$+(a_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$ are sixth

augmentation coefficient for category 1, 2 and 3
$+(a_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)$ are Seventh augmentation coefficient for category 1, 2 and 3
$+(a_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)$ $+(a_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)$ $+(a_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)$ are eighth augmentation coefficient for 1,2,3
$+(a_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{44}}{dt} =$$
$$(b_{44})^{(9)}T_{45} -$$
$$\begin{bmatrix} (b_{44}^{'})^{(9)}\boxed{-(b_{44}^{''})^{(9)}(G_{47},t)} & \boxed{-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{bmatrix} T_{13}$$

$$\frac{dT_{45}}{dt} =$$
$$(b_{45})^{(9)}T_{44} - \begin{bmatrix} (b_{45}^{'})^{(9)}\boxed{-(b_{45}^{''})^{(9)}(G_{47},t)} & \boxed{-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{bmatrix} T_{14}$$

$$\frac{dT_{46}}{dt} =$$
$$(b_{46})^{(9)}T_{45} - \begin{bmatrix} (b_{46}^{'})^{(9)}\boxed{-(b_{46}^{''})^{(9)}(G_{47},t)} & \boxed{-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{bmatrix} T_{15}$$

Where $-(b_{44}^{''})^{(9)}(G_{47},t)$, $-(b_{45}^{''})^{(9)}(G_{47},t)$, $-(b_{46}^{''})^{(9)}(G_{47},t)$ are first detrition coefficients for category 1, 2 and 3
$-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3
$-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3
$-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3
$-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3
$-(b_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)$ are sixth detrition coefficients for category 1, 2 and 3
$-(b_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)$ are seventh detrition coefficients for category 1, 2 and 3

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| $\boxed{-(b_{37}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{36}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{38}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$ are eighth detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$ are ninth detrition coefficients for category 1, 2 and 3 | |
| Where we suppose | |
| $(a_i)^{(1)}, (a_i^{'})^{(1)}, (a_i^{''})^{(1)}, (b_i)^{(1)}, (b_i^{'})^{(1)}, (b_i^{''})^{(1)} > 0, \quad i,j = 13,14,15$<br><br>The functions $(a_i^{''})^{(1)}, (b_i^{''})^{(1)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(1)}$, $(r_i)^{(1)}$:<br><br>$(a_i^{''})^{(1)}(T_{14},t) \leq (p_i)^{(1)} \leq (\hat{A}_{13})^{(1)}$<br><br>$(b_i^{''})^{(1)}(G,t) \leq (r_i)^{(1)} \leq (b_i^{'})^{(1)} \leq (\hat{B}_{13})^{(1)}$ | *97* |
| $lim_{T_2\to\infty}(a_i^{''})^{(1)}(T_{14},t) = (p_i)^{(1)}$<br><br>$\lim_{G\to\infty}(b_i^{''})^{(1)}(G,t) = (r_i)^{(1)}$<br><br>**Definition of** $(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}$ :<br><br>Where $\boxed{(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}}$ are positive constants and $\boxed{i = 13,14,15}$ | 98 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(1)}(T_{14}^{'},t) - (a_i^{''})^{(1)}(T_{14},t)\rvert \leq (\hat{k}_{13})^{(1)}\lvert T_{14} - T_{14}^{'}\rvert e^{-(\hat{M}_{13})^{(1)}t}$<br><br>$\lvert(b_i^{''})^{(1)}(G^{'},t) - (b_i^{''})^{(1)}(G,t)\rvert < (\hat{k}_{13})^{(1)}\lVert G - G^{'}\rVert e^{-(\hat{M}_{13})^{(1)}t}$ | 99 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(1)}(T_{14}^{'},t)$ and $(a_i^{''})^{(1)}(T_{14},t)$ . $(T_{14}^{'},t)$ and $(T_{14},t)$ are points belonging to the interval $\left[(\hat{k}_{13})^{(1)}, (\hat{M}_{13})^{(1)}\right]$. It is to be noted that $(a_i^{''})^{(1)}(T_{14},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{13})^{(1)} = 1$ then the function $(a_i^{''})^{(1)}(T_{14},t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}$ :<br><br>$(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}$, are positive constants<br><br>$\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}} < 1$ | 100 |
| **Definition of** $(\hat{P}_{13})^{(1)}, (\hat{Q}_{13})^{(1)}$ :<br><br>There exists two constants $(\hat{P}_{13})^{(1)}$ and $(\hat{Q}_{13})^{(1)}$ which together With $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}, (\hat{A}_{13})^{(1)}$ and $(\hat{B}_{13})^{(1)}$ and the constants $(a_i)^{(1)}, (a_i^{'})^{(1)}, (b_i)^{(1)}, (b_i^{'})^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}, i = 13,14,15,$ | 101 |

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| satisfy the inequalities<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(a_i)^{(1)} + (a_i')^{(1)} + (\hat{A}_{13})^{(1)} + (\hat{P}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(b_i)^{(1)} + (b_i')^{(1)} + (\hat{B}_{13})^{(1)} + (\hat{Q}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$ | |
| Where we suppose | |
| $(a_i)^{(2)}, (a_i')^{(2)}, (a_i'')^{(2)}, (b_i)^{(2)}, (b_i')^{(2)}, (b_i'')^{(2)} > 0, \quad i,j = 16,17,18$ | |
| The functions $(a_i'')^{(2)}, (b_i'')^{(2)}$ are positive continuous increasing and bounded. | |
| **Definition of** $(p_i)^{(2)}$, $(r_i)^{(2)}$: | |
| $(a_i'')^{(2)}(T_{17}, t) \leq (p_i)^{(2)} \leq (\hat{A}_{16})^{(2)}$ | 102 |
| $(b_i'')^{(2)}(G_{19}, t) \leq (r_i)^{(2)} \leq (b_i')^{(2)} \leq (\hat{B}_{16})^{(2)}$ | 103 |
| $\lim_{T_2 \to \infty} (a_i'')^{(2)}(T_{17}, t) = (p_i)^{(2)}$ | 104 |
| $\lim_{G \to \infty} (b_i'')^{(2)}((G_{19}), t) = (r_i)^{(2)}$ | 105 |
| **Definition of** $(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}$ :<br>Where $(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}, (p_i)^{(2)}, (r_i)^{(2)}$ are positive constants and $i = 16,17,18$ | 106 |
| They satisfy Lipschitz condition: | |
| $\lvert(a_i'')^{(2)}(T_{17}', t) - (a_i'')^{(2)}(T_{17}, t)\rvert \leq (\hat{k}_{16})^{(2)} \lvert T_{17} - T_{17}' \rvert e^{-(\hat{M}_{16})^{(2)}t}$ | 107 |
| $\lvert(b_i'')^{(2)}((G_{19})', t) - (b_i'')^{(2)}((G_{19}), t)\rvert < (\hat{k}_{16})^{(2)} \lvert\lvert(G_{19}) - (G_{19})'\rvert\rvert e^{-(\hat{M}_{16})^{(2)}t}$ | 108 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(2)}(T_{17}', t)$ and $(a_i'')^{(2)}(T_{17}, t)$ . $(T_{17}', t)$ and $(T_{17}, t)$ are points belonging to the interval $[(\hat{k}_{16})^{(2)}, (\hat{M}_{16})^{(2)}]$ . It is to be noted that $(a_i'')^{(2)}(T_{17}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{16})^{(2)} = 1$ then the function $(a_i'')^{(2)}(T_{17}, t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$ : | |
| $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$, are positive constants<br>$\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}, \frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ | 109 |
| **Definition of** $(\hat{P}_{13})^{(2)}, (\hat{Q}_{13})^{(2)}$ :<br>There exists two constants $(\hat{P}_{16})^{(2)}$ and $(\hat{Q}_{16})^{(2)}$ which together | |

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| with $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}, (\hat{A}_{16})^{(2)} and\ (\hat{B}_{16})^{(2)}$ and the constants $(a_i)^{(2)}, (a_i^{'})^{(2)}, (b_i)^{(2)}, (b_i^{'})^{(2)}, (p_i)^{(2)}, \ (r_i)^{(2)}, i = 16,17,18,$ satisfy the inequalities | |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(a_i)^{(2)} + (a_i^{'})^{(2)} + (\hat{A}_{16})^{(2)} + (\hat{P}_{16})^{(2)} (\hat{k}_{16})^{(2)}] < 1$ | 110 |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(b_i)^{(2)} + (b_i^{'})^{(2)} + (\hat{B}_{16})^{(2)} + (\hat{Q}_{16})^{(2)} (\hat{k}_{16})^{(2)}] < 1$ | 111 |
| Where we suppose | |
| $(a_i)^{(3)}, (a_i^{'})^{(3)}, (a_i^{''})^{(3)}, (b_i)^{(3)}, (b_i^{'})^{(3)}, (b_i^{''})^{(3)} > 0, \quad i, j = 20,21,22$<br>The functions $(a_i^{''})^{(3)}, (b_i^{''})^{(3)}$ are positive continuous increasing and bounded.<br>**Definition of** $(p_i)^{(3)}$, $(r_i)^{(3)}$:<br>$(a_i^{''})^{(3)}(T_{21}, t) \leq (p_i)^{(3)} \leq (\hat{A}_{20})^{(3)}$<br>$(b_i^{''})^{(3)}(G_{23}, t) \leq (r_i)^{(3)} \leq (b_i^{'})^{(3)} \leq (\hat{B}_{20})^{(3)}$ | 112 |
| $lim_{T_2 \to \infty}(a_i^{''})^{(3)}(T_{21}, t) = (p_i)^{(3)}$<br>$\lim_{G \to \infty}(b_i^{''})^{(3)}(G_{23}, t) = (r_i)^{(3)}$<br>**Definition of** $(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}$ :<br>Where $\boxed{(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}}$ are positive constants and $\boxed{i = 20,21,22}$ | 113 |
| They satisfy Lipschitz condition:<br>$\lvert(a_i^{''})^{(3)}(T_{21}^{'}, t) - (a_i^{''})^{(3)}(T_{21}, t)\rvert \leq (\hat{k}_{20})^{(3)} \lvert T_{21} - T_{21}^{'}\rvert e^{-(\hat{M}_{20})^{(3)}t}$<br>$\lvert(b_i^{''})^{(3)}(G_{23}^{'}, t) - (b_i^{''})^{(3)}(G_{23}, t)\rvert < (\hat{k}_{20})^{(3)} \lvert\lvert G_{23} - G_{23}^{'}\rvert\rvert e^{-(\hat{M}_{20})^{(3)}t}$ | 114 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(3)}(T_{21}^{'}, t)$ and $(a_i^{''})^{(3)}(T_{21}, t)$ . $(T_{21}^{'}, t)$ And $(T_{21}, t)$ are points belonging to the interval $[(\hat{k}_{20})^{(3)}, (\hat{M}_{20})^{(3)}]$. It is to be noted that $(a_i^{''})^{(3)}(T_{21}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{20})^{(3)} = 1$ then the function $(a_i^{''})^{(3)}(T_{21}, t)$, the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$ :<br>$(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$, are positive constants<br>$\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}, \frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}} < 1$ | 115 |

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| There exists two constants There exists two constants $(\hat{P}_{20})^{(3)}$ and $(\hat{Q}_{20})^{(3)}$ which together with $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}, (\hat{A}_{20})^{(3)} and\ (\hat{B}_{20})^{(3)}$ and the constants $(a_i)^{(3)}, (a_i^{'})^{(3)}, (b_i)^{(3)}, (b_i^{'})^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}, i = 20,21,22,$ satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(a_i)^{(3)} + (a_i^{'})^{(3)} + (\hat{A}_{20})^{(3)} + (\hat{P}_{20})^{(3)} (\hat{k}_{20})^{(3)}] < 1$<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(b_i)^{(3)} + (b_i^{'})^{(3)} + (\hat{B}_{20})^{(3)} + (\hat{Q}_{20})^{(3)} (\hat{k}_{20})^{(3)}] < 1$ | 116 |
| Where we suppose | |
| $(a_i)^{(4)}, (a_i^{'})^{(4)}, (a_i^{''})^{(4)}, (b_i)^{(4)}, (b_i^{'})^{(4)}, (b_i^{''})^{(4)} > 0, \quad i,j = 24,25,26$<br><br>The functions $(a_i^{''})^{(4)}, (b_i^{''})^{(4)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(4)}$, $(r_i)^{(4)}$:<br><br>$(a_i^{''})^{(4)}(T_{25}, t) \leq (p_i)^{(4)} \leq (\hat{A}_{24})^{(4)}$<br><br>$(b_i^{''})^{(4)}((G_{27}), t) \leq (r_i)^{(4)} \leq (b_i^{'})^{(4)} \leq (\hat{B}_{24})^{(4)}$ | 117 |
| $lim_{T_2 \to \infty} (a_i^{''})^{(4)}(T_{25}, t) = (p_i)^{(4)}$<br><br>$\lim_{G \to \infty} (b_i^{''})^{(4)}((G_{27}), t) = (r_i)^{(4)}$<br><br>**Definition of** $(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}$ :<br><br>Where $\boxed{(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}, (p_i)^{(4)}, (r_i)^{(4)}}$ are positive constants and $\boxed{i = 24,25,26}$ | 118 |
| They satisfy Lipschitz condition:<br><br>$\|(a_i^{''})^{(4)}(T_{25}^{'}, t) - (a_i^{''})^{(4)}(T_{25}, t)\| \leq (\hat{k}_{24})^{(4)} \|T_{25} - T_{25}^{'}\| e^{-(\hat{M}_{24})^{(4)} t}$<br><br>$\|(b_i^{''})^{(4)}((G_{27})', t) - (b_i^{''})^{(4)}((G_{27}), t)\| < (\hat{k}_{24})^{(4)} \|(G_{27}) - (G_{27})'\| e^{-(\hat{M}_{24})^{(4)} t}$ | 119 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(4)}(T_{25}^{'}, t)$ and $(a_i^{''})^{(4)}(T_{25}, t)$ . $(T_{25}^{'}, t)$ and $(T_{25}, t)$ are points belonging to the interval $[(\hat{k}_{24})^{(4)}, (\hat{M}_{24})^{(4)}]$. It is to be noted that $(a_i^{''})^{(4)}(T_{25}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{24})^{(4)} = 1$ then the function $(a_i^{''})^{(4)}(T_{25}, t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$ :<br><br>$(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$, are positive constants | 120 |

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| $\frac{(a_i)^{(4)}}{(\hat{M}_{24})^{(4)}}, \frac{(b_i)^{(4)}}{(\hat{M}_{24})^{(4)}} < 1$ | |
| **Definition of** $(\hat{P}_{24})^{(4)}, (\hat{Q}_{24})^{(4)}$ :<br><br>There exists two constants $(\hat{P}_{24})^{(4)}$ and $(\hat{Q}_{24})^{(4)}$ which together with $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}, (\hat{A}_{24})^{(4)} and\ (\hat{B}_{24})^{(4)}$ and the constants $(a_i)^{(4)}, (a_i^{'})^{(4)}, (b_i)^{(4)}, (b_i^{'})^{(4)}, (p_i)^{(4)}, (r_i)^{(4)}, i = 24,25,26$, satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{24})^{(4)}}[(a_i)^{(4)} + (a_i^{'})^{(4)} + (\hat{A}_{24})^{(4)} + (\hat{P}_{24})^{(4)} (\hat{k}_{24})^{(4)}] < 1$<br><br>$\frac{1}{(\hat{M}_{24})^{(4)}}[(b_i)^{(4)} + (b_i^{'})^{(4)} + (\hat{B}_{24})^{(4)} + (\hat{Q}_{24})^{(4)} (\hat{k}_{24})^{(4)}] < 1$ | 121 |
| Where we suppose | |
| $(a_i)^{(5)}, (a_i^{'})^{(5)}, (a_i^{''})^{(5)}, (b_i)^{(5)}, (b_i^{'})^{(5)}, (b_i^{''})^{(5)} > 0, \quad i,j = 28,29,30$<br><br>The functions $(a_i^{''})^{(5)}, (b_i^{''})^{(5)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(5)}, (r_i)^{(5)}$:<br><br>$(a_i^{''})^{(5)}(T_{29}, t) \leq (p_i)^{(5)} \leq (\hat{A}_{28})^{(5)}$<br><br>$(b_i^{''})^{(5)}((G_{31}), t) \leq (r_i)^{(5)} \leq (b_i^{'})^{(5)} \leq (\hat{B}_{28})^{(5)}$ | 122 |
| $lim_{T_2 \to \infty} (a_i^{''})^{(5)} (T_{29}, t) = (p_i)^{(5)}$<br><br>$\lim_{G \to \infty} (b_i^{''})^{(5)} (G_{31}, t) = (r_i)^{(5)}$<br><br>**Definition of** $(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}$ :<br><br>Where $\boxed{(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}, (p_i)^{(5)}, (r_i)^{(5)}}$ are positive constants and $\boxed{i = 28,29,30}$ | 123 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(5)}(T_{29}^{'}, t) - (a_i^{''})^{(5)}(T_{29}, t)\rvert \leq (\hat{k}_{28})^{(5)} \lvert T_{29} - T_{29}^{'}\rvert e^{-(\hat{M}_{28})^{(5)} t}$<br><br>$\lvert(b_i^{''})^{(5)}((G_{31})', t) - (b_i^{''})^{(5)}((G_{31}), t)\rvert < (\hat{k}_{28})^{(5)} \lvert\lvert(G_{31}) - (G_{31})'\rvert\rvert e^{-(\hat{M}_{28})^{(5)} t}$ | 124 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(5)}(T_{29}^{'}, t)$ and $(a_i^{''})^{(5)}(T_{29}, t)$ . $(T_{29}^{'}, t)$ and $(T_{29}, t)$ are points belonging to the interval $[(\hat{k}_{28})^{(5)}, (\hat{M}_{28})^{(5)}]$. It is to be noted that $(a_i^{''})^{(5)}(T_{29}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{28})^{(5)} = 1$ then the function $(a_i^{''})^{(5)}(T_{29}, t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}$ : | 125 |

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| $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}$, are positive constants<br><br>$\frac{(a_i)^{(5)}}{(\hat{M}_{28})^{(5)}}$ , $\frac{(b_i)^{(5)}}{(\hat{M}_{28})^{(5)}} < 1$ | |
| **Definition of** $(\hat{P}_{28})^{(5)}, (\hat{Q}_{28})^{(5)}$ :<br><br>There exists two constants $(\hat{P}_{28})^{(5)}$ and $(\hat{Q}_{28})^{(5)}$ which together with $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}, (\hat{A}_{28})^{(5)} and\ (\hat{B}_{28})^{(5)}$ and the constants $(a_i)^{(5)}, (a_i^{'})^{(5)}, (b_i)^{(5)}, (b_i^{'})^{(5)}, (p_i)^{(5)}, (r_i)^{(5)}, i = 28,29,30,$ satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{28})^{(5)}}[(a_i)^{(5)} + (a_i^{'})^{(5)} + (\hat{A}_{28})^{(5)} + (\hat{P}_{28})^{(5)} (\hat{k}_{28})^{(5)}] < 1$<br><br>$\frac{1}{(\hat{M}_{28})^{(5)}}[(b_i)^{(5)} + (b_i^{'})^{(5)} + (\hat{B}_{28})^{(5)} + (\hat{Q}_{28})^{(5)} (\hat{k}_{28})^{(5)}] < 1$ | 126 |
| Where we suppose | |
| $(a_i)^{(6)}, (a_i^{'})^{(6)}, (a_i^{''})^{(6)}, (b_i)^{(6)}, (b_i^{'})^{(6)}, (b_i^{''})^{(6)} > 0, \quad i,j = 32,33,34$<br><br>The functions $(a_i^{''})^{(6)}, (b_i^{''})^{(6)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(6)}$, $(r_i)^{(6)}$:<br><br>$(a_i^{''})^{(6)}(T_{33}, t) \leq (p_i)^{(6)} \leq (\hat{A}_{32})^{(6)}$<br><br>$(b_i^{''})^{(6)}((G_{35}), t) \leq (r_i)^{(6)} \leq (b_i^{'})^{(6)} \leq (\hat{B}_{32})^{(6)}$ | 127 |
| $lim_{T_2 \to \infty} (a_i^{''})^{(6)} (T_{33}, t) = (p_i)^{(6)}$<br><br>$\lim_{G \to \infty} (b_i^{''})^{(6)} ((G_{35}), t) = (r_i)^{(6)}$<br><br>**Definition of** $(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}$ :<br><br>Where $\boxed{(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}, (p_i)^{(6)}, (r_i)^{(6)}}$ are positive constants and $\boxed{i = 32,33,34}$ | 128 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(6)}(T_{33}^{'}, t) - (a_i^{''})^{(6)}(T_{33}, t)\rvert \leq (\hat{k}_{32})^{(6)} \lvert T_{33} - T_{33}^{'}\rvert e^{-(\hat{M}_{32})^{(6)} t}$<br><br>$\lvert(b_i^{''})^{(6)}((G_{35})', t) - (b_i^{''})^{(6)}((G_{35}), t)\rvert < (\hat{k}_{32})^{(6)} \lvert\lvert(G_{35}) - (G_{35})'\rvert\rvert e^{-(\hat{M}_{32})^{(6)} t}$ | |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(6)}(T_{33}^{'}, t)$ and $(a_i^{''})^{(6)}(T_{33}, t)$ . $(T_{33}^{'}, t)$ and $(T_{33}, t)$ are points belonging to the interval $[(\hat{k}_{32})^{(6)}, (\hat{M}_{32})^{(6)}]$. It is to be noted that $(a_i^{''})^{(6)}(T_{33}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{32})^{(6)} = 1$ then the function $(a_i^{''})^{(6)}(T_{33}, t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |

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| **Definition of** $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}$ :<br><br>$(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)},$ are positive constants<br><br>$\frac{(a_i)^{(6)}}{(\hat{M}_{32})^{(6)}}$ , $\frac{(b_i)^{(6)}}{(\hat{M}_{32})^{(6)}} < 1$ | 129 |
| **Definition of** $(\hat{P}_{32})^{(6)}, (\hat{Q}_{32})^{(6)}$ :<br><br>There exists two constants $(\hat{P}_{32})^{(6)}$ and $(\hat{Q}_{32})^{(6)}$ which together with $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}, (\hat{A}_{32})^{(6)} and\ (\hat{B}_{32})^{(6)}$ and the constants $(a_i)^{(6)}, (a_i')^{(6)}, (b_i)^{(6)}, (b_i')^{(6)}, (p_i)^{(6)}, (r_i)^{(6)}, i = 32,33,34,$ satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{32})^{(6)}}[(a_i)^{(6)} + (a_i')^{(6)} + (\hat{A}_{32})^{(6)} + (\hat{P}_{32})^{(6)} (\hat{k}_{32})^{(6)}] < 1$<br><br>$\frac{1}{(\hat{M}_{32})^{(6)}}[(b_i)^{(6)} + (b_i')^{(6)} + (\hat{B}_{32})^{(6)} + (\hat{Q}_{32})^{(6)} (\hat{k}_{32})^{(6)}] < 1$ | 130 |
| Where we suppose | |
| (EE) $(a_i)^{(7)}, (a_i')^{(7)}, (a_i'')^{(7)}, (b_i)^{(7)}, (b_i')^{(7)}, (b_i'')^{(7)} > 0, \quad i, j = 36,37,38$<br><br>(FF) The functions $(a_i'')^{(7)}, (b_i'')^{(7)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(7)}$, $(r_i)^{(7)}$:<br><br>$(a_i'')^{(7)}(T_{37}, t) \le (p_i)^{(7)} \le (\hat{A}_{36})^{(7)}$<br><br>$(b_i'')^{(7)}(G_{39}, t) \le (r_i)^{(7)} \le (b_i')^{(7)} \le (\hat{B}_{36})^{(7)}$ | 131 |
| *(GG)* $\lim_{T_2 \to \infty} (a_i'')^{(7)}(T_{37}, t) = (p_i)^{(7)}$<br>*(HH)*<br>$\lim_{G \to \infty} (b_i'')^{(7)}((G_{39}), t) = (r_i)^{(7)}$<br><br>**Definition of** $(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}$ :<br><br>Where $\boxed{(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}, (p_i)^{(7)}, (r_i)^{(7)}}$ are positive constants and $\boxed{i = 36,37,38}$ | 132 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i'')^{(7)}(T_{37}', t) - (a_i'')^{(7)}(T_{37}, t)\rvert \le (\hat{k}_{36})^{(7)} \lvert T_{37} - T_{37}' \rvert e^{-(\hat{M}_{36})^{(7)} t}$<br><br>$\lvert(b_i'')^{(7)}((G_{39})', t) - (b_i'')^{(7)}((G_{39}), t)\rvert < (\hat{k}_{36})^{(7)} \lvert\lvert(G_{39}) - (G_{39})'\rvert\rvert e^{-(\hat{M}_{36})^{(7)} t}$ | 133 |

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| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(7)}(T_{37}^{'},t)$ and$(a_i^{''})^{(7)}(T_{37},t)$ . $(T_{37}^{'},t)$ and $(T_{37},t)$ are points belonging to the interval $\left[(\hat{k}_{36})^{(7)},(\widehat{M}_{36})^{(7)}\right]$. It is to be noted that $(a_i^{''})^{(7)}(T_{37},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\widehat{M}_{36})^{(7)}=1$ then the function $(a_i^{''})^{(7)}(T_{37},t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\widehat{M}_{36})^{(7)},(\hat{k}_{36})^{(7)}$ : <br><br>(II) $(\widehat{M}_{36})^{(7)},(\hat{k}_{36})^{(7)},$ are positive constants <br><br> $\frac{(a_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}$ , $\frac{(b_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}<1$ | 134 |
| **Definition of** $(\hat{P}_{36})^{(7)},(\hat{Q}_{36})^{(7)}$ : <br><br>(JJ) There exists two constants $(\hat{P}_{36})^{(7)}$ and $(\hat{Q}_{36})^{(7)}$ which together with $(\widehat{M}_{36})^{(7)},(\hat{k}_{36})^{(7)},(\hat{A}_{36})^{(7)} and\ (\hat{B}_{36})^{(7)}$ and the constants $(a_i)^{(7)},(a_i^{'})^{(7)},(b_i)^{(7)},(b_i^{'})^{(7)},(p_i)^{(7)},\ (r_i)^{(7)},i=36,37,38$, satisfy the inequalities <br><br> $\frac{1}{(\widehat{M}_{36})^{(7)}}[(a_i)^{(7)}+(a_i^{'})^{(7)}+(\hat{A}_{36})^{(7)}+(\hat{P}_{36})^{(7)}(\hat{k}_{36})^{(7)}]<1$ <br><br> $\frac{1}{(\widehat{M}_{36})^{(7)}}[(b_i)^{(7)}+(b_i^{'})^{(7)}+(\hat{B}_{36})^{(7)}+(\hat{Q}_{36})^{(7)}(\hat{k}_{36})^{(7)}]<1$ | 135 |
| Where we suppose | |
| $(a_i)^{(8)},(a_i^{'})^{(8)},(a_i^{''})^{(8)},(b_i)^{(8)},(b_i^{'})^{(8)},(b_i^{''})^{(8)}>0,\quad i,j=40,41,42$ | 136 |
| The functions $(a_i^{''})^{(8)},(b_i^{''})^{(8)}$ are positive continuous increasing and bounded | |
| **Definition of** $(p_i)^{(8)},\ (r_i)^{(8)}$: | 137 |
| $(a_i^{''})^{(8)}(T_{41},t)\leq(p_i)^{(8)}\leq(\hat{A}_{40})^{(8)}$ | 138 |
| $(b_i^{''})^{(8)}((G_{43}),t)\leq\ (r_i)^{(8)}\leq(b_i^{'})^{(8)}\leq(\hat{B}_{40})^{(8)}$ | 139 |
| $\lim_{T_2\to\infty}(a_i^{''})^{(8)}(T_{41},t)=(p_i)^{(8)}$ | 140 |
| $\lim_{G\to\infty}(b_i^{''})^{(8)}\left((G_{43}),t\right)=\ (r_i)^{(8)}$ | 141 |
| **Definition of** $(\hat{A}_{40})^{(8)},(\hat{B}_{40})^{(8)}$ : <br><br> Where $\boxed{(\hat{A}_{40})^{(8)},(\hat{B}_{40})^{(8)},(p_i)^{(8)},\ (r_i)^{(8)}}$ are positive constants and $\boxed{i=40,41,42}$ | |

They satisfy Lipschitz condition:

$$|(a_i'')^{(8)}(T_{41}', t) - (a_i'')^{(8)}(T_{41}, t)| \leq (\hat{k}_{40})^{(8)}|T_{41} - T_{41}'|e^{-(\hat{M}_{40})^{(8)}t} \quad (142)$$

$$|(b_i'')^{(8)}((G_{43})', t) - (b_i'')^{(8)}((G_{43}), t)| < (\hat{k}_{40})^{(8)}||(G_{43}) - (G_{43})'||e^{-(\hat{M}_{40})^{(8)}t} \quad (143)$$

With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(8)}(T_{41}', t)$ and $(a_i'')^{(8)}(T_{41}, t)$ . $(T_{41}', t)$ and $(T_{41}, t)$ are points belonging to the interval $[(\hat{k}_{40})^{(8)}, (\hat{M}_{40})^{(8)}]$. It is to be noted that $(a_i'')^{(8)}(T_{41}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{40})^{(8)} = 1$ then the function $(a_i'')^{(8)}(T_{41}, t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

**Definition of** $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}$ :

$(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}$, are positive constants

$$\frac{(a_i)^{(8)}}{(\hat{M}_{40})^{(8)}}, \frac{(b_i)^{(8)}}{(\hat{M}_{40})^{(8)}} < 1 \quad (144)$$

**Definition of** $(\hat{P}_{40})^{(8)}, (\hat{Q}_{40})^{(8)}$ :

There exists two constants $(\hat{P}_{40})^{(8)}$ and $(\hat{Q}_{40})^{(8)}$ which together with $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}, (\hat{A}_{40})^{(8)}$ $(\hat{B}_{40})^{(8)}$ and the constants $(a_i)^{(8)}, (a_i')^{(8)}, (b_i)^{(8)}, (b_i')^{(8)}, (p_i)^{(8)}, (r_i)^{(8)}, i = 40,41,42,$
Satisfy the inequalities

$$\frac{1}{(\hat{M}_{40})^{(8)}}[(a_i)^{(8)} + (a_i')^{(8)} + (\hat{A}_{40})^{(8)} + (\hat{P}_{40})^{(8)}(\hat{k}_{40})^{(8)}] < 1 \quad (145)$$

$$\frac{1}{(\hat{M}_{40})^{(8)}}[(b_i)^{(8)} + (b_i')^{(8)} + (\hat{B}_{40})^{(8)} + (\hat{Q}_{40})^{(8)}(\hat{k}_{40})^{(8)}] < 1 \quad (146)$$

Where we suppose

$$(a_i)^{(9)}, (a_i')^{(9)}, (a_i'')^{(9)}, (b_i)^{(9)}, (b_i')^{(9)}, (b_i'')^{(9)} > 0, \quad i, j = 44,45,46 \quad (146A)$$

The functions $(a_i'')^{(9)}, (b_i'')^{(9)}$ are positive continuous increasing and bounded.

**Definition of** $(p_i)^{(9)}$, $(r_i)^{(9)}$:

$$(a_i'')^{(9)}(T_{45}, t) \leq (p_i)^{(9)} \leq (\hat{A}_{44})^{(9)}$$

$$(b_i'')^{(9)}(G_{47}, t) \leq (r_i)^{(9)} \leq (b_i')^{(9)} \leq (\hat{B}_{44})^{(9)}$$

$$\lim_{T_2 \to \infty} (a_i'')^{(9)}(T_{45}, t) = (p_i)^{(9)}$$

$$\lim_{G \to \infty} (b_i'')^{(9)}(G_{47}, t) = (r_i)^{(9)}$$

**Definition of** $(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}$ :

Where $\boxed{(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}, (p_i)^{(9)}, (r_i)^{(9)}}$ are positive constants and $\boxed{i = 44,45,46}$

They satisfy Lipschitz condition:

$$|(a_i'')^{(9)}(T_{45}', t) - (a_i'')^{(9)}(T_{45}, t)| \leq (\hat{k}_{44})^{(9)}|T_{45} - T_{45}'|e^{-(\hat{M}_{44})^{(9)}t}$$

$$|(b_i'')^{(9)}((G_{47})', t) - (b_i'')^{(9)}\big((G_{47}), t\big)| < (\hat{k}_{44})^{(9)}||(G_{47}) - (G_{47})'||e^{-(\hat{M}_{44})^{(9)}t}$$

With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(9)}(T_{45}', t)$ and $(a_i'')^{(9)}(T_{45}, t)$ . $(T_{45}', t)$ and $(T_{45}, t)$ are points belonging to the interval $\left[(\hat{k}_{44})^{(9)}, (\hat{M}_{44})^{(9)}\right]$. It is to be noted that $(a_i'')^{(9)}(T_{45}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{44})^{(9)} = 1$ then the function $(a_i'')^{(9)}(T_{45}, t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

**Definition of** $(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}$ :

$(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}$, are positive constants

$$\frac{(a_i)^{(9)}}{(\hat{M}_{44})^{(9)}} \ , \frac{(b_i)^{(9)}}{(\hat{M}_{44})^{(9)}} < 1$$

**Definition of** $(\hat{P}_{44})^{(9)}, (\hat{Q}_{44})^{(9)}$ :

There exists two constants $(\hat{P}_{44})^{(9)}$ and $(\hat{Q}_{44})^{(9)}$ which together with $(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}, (\hat{A}_{44})^{(9)} and\ (\hat{B}_{44})^{(9)}$ and the constants $(a_i)^{(9)}, (a_i')^{(9)}, (b_i)^{(9)}, (b_i')^{(9)}, (p_i)^{(9)}, (r_i)^{(9)}, i = 44,45,46,$ satisfy the inequalities

$$\frac{1}{(\hat{M}_{44})^{(9)}}[(a_i)^{(9)} + (a_i')^{(9)} + (\hat{A}_{44})^{(9)} + (\hat{P}_{44})^{(9)}(\hat{k}_{44})^{(9)}] < 1$$

$$\frac{1}{(\hat{M}_{44})^{(9)}}[(b_i)^{(9)} + (b_i')^{(9)} + (\hat{B}_{44})^{(9)} + (\hat{Q}_{44})^{(9)}(\hat{k}_{44})^{(9)}] < 1$$

**Theorem 1:** if the conditions above are fulfilled, there exists a solution satisfying the conditions 147

Definition of $G_i(0), T_i(0)$ :

$$G_i(t) \leq (\hat{P}_{13})^{(1)}e^{(\hat{M}_{13})^{(1)}t} \ , \ \boxed{G_i(0) = G_i^0 > 0}$$

$$T_i(t) \leq (\hat{Q}_{13})^{(1)}e^{(\hat{M}_{13})^{(1)}t} \ , \ \boxed{T_i(0) = T_i^0 > 0}$$

**Theorem 2 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions 148

**Definition of** $G_i(0), T_i(0)$

$$G_i(t) \leq (\hat{P}_{16})^{(2)}e^{(\hat{M}_{16})^{(2)}t} \ , \ G_i(0) = G_i^0 > 0$$

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| $T_i(t) \le (\hat{Q}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ , $T_i(0) = T_i^0 > 0$ | |
| **Theorem 3 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>$G_i(t) \le (\hat{P}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ , $G_i(0) = G_i^0 > 0$<br>$T_i(t) \le (\hat{Q}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ , $T_i(0) = T_i^0 > 0$ | 149 |
| **Theorem 4 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \le (\hat{P}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le (\hat{Q}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 150 |
| **Theorem 5 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \le (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le (\hat{Q}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 151 |
| **Theorem 6 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \le (\hat{P}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le (\hat{Q}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 152 |
| **Theorem 7:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \le (\hat{P}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le (\hat{Q}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 |
| **Theorem 8:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \le (\hat{P}_{40})^{(8)} e^{(\hat{M}_{40})^{(8)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$ | 153 A |

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| $T_i(t) \leq (\hat{Q}_{40})^{(8)} e^{(\hat{M}_{40})^{(8)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | |
| **Theorem 9:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br><br>**Definition of** $G_i(0), T_i(0)$ :<br><br>$G_i(t) \leq (\hat{P}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br><br>$T_i(t) \leq (\hat{Q}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 B |
| **Proof:** Consider operator $\mathcal{A}^{(1)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 154 |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{13})^{(1)}$ , $T_i^0 \leq (\hat{Q}_{13})^{(1)}$, | 155 |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ | 156 |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ | 157 |
| By<br><br>$\bar{G}_{13}(t) = G_{13}^0 + \int_0^t \left[ (a_{13})^{(1)} G_{14}(s_{(13)}) - \left( (a_{13}')^{(1)} + a_{13}'')^{(1)} (T_{14}(s_{(13)}), s_{(13)}) \right) G_{13}(s_{(13)}) \right] ds_{(13)}$ | 158 |
| $\bar{G}_{14}(t) = G_{14}^0 + \int_0^t \left[ (a_{14})^{(1)} G_{13}(s_{(13)}) - \left( (a_{14}')^{(1)} + (a_{14}'')^{(1)} (T_{14}(s_{(13)}), s_{(13)}) \right) G_{14}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{G}_{15}(t) = G_{15}^0 + \int_0^t \left[ (a_{15})^{(1)} G_{14}(s_{(13)}) - \left( (a_{15}')^{(1)} + (a_{15}'')^{(1)} (T_{14}(s_{(13)}), s_{(13)}) \right) G_{15}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{T}_{13}(t) = T_{13}^0 + \int_0^t \left[ (b_{13})^{(1)} T_{14}(s_{(13)}) - \left( (b_{13}')^{(1)} - (b_{13}'')^{(1)} (G(s_{(13)}), s_{(13)}) \right) T_{13}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{T}_{14}(t) = T_{14}^0 + \int_0^t \left[ (b_{14})^{(1)} T_{13}(s_{(13)}) - \left( (b_{14}')^{(1)} - (b_{14}'')^{(1)} (G(s_{(13)}), s_{(13)}) \right) T_{14}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{T}_{15}(t) = T_{15}^0 + \int_0^t \left[ (b_{15})^{(1)} T_{14}(s_{(13)}) - \left( (b_{15}')^{(1)} - (b_{15}'')^{(1)} (G(s_{(13)}), s_{(13)}) \right) T_{15}(s_{(13)}) \right] ds_{(13)}$<br><br>Where $s_{(13)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br><br>Consider operator $\mathcal{A}^{(2)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 159 |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{16})^{(2)}$ , $T_i^0 \leq (\hat{Q}_{16})^{(2)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ | |

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| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{16})^{(2)} e^{(\bar{M}_{16})^{(2)} t}$ | |
| By<br>$\bar{G}_{16}(t) = G_{16}^0 + \int_0^t \left[ (a_{16})^{(2)} G_{17}(s_{(16)}) - \left( (a_{16}')^{(2)} + a_{16}'')^{(2)} (T_{17}(s_{(16)}), s_{(16)}) \right) G_{16}(s_{(16)}) \right] ds_{(16)}$ | 160 |
| $\bar{G}_{17}(t) = G_{17}^0 + \int_0^t \left[ (a_{17})^{(2)} G_{16}(s_{(16)}) - \left( (a_{17}')^{(2)} + (a_{17}'')^{(2)} (T_{17}(s_{(16)}), s_{(17)}) \right) G_{17}(s_{(16)}) \right] ds_{(16)}$ | |
| $\bar{G}_{18}(t) = G_{18}^0 + \int_0^t \left[ (a_{18})^{(2)} G_{17}(s_{(16)}) - \left( (a_{18}')^{(2)} + (a_{18}'')^{(2)} (T_{17}(s_{(16)}), s_{(16)}) \right) G_{18}(s_{(16)}) \right] ds_{(16)}$ | |
| $\bar{T}_{16}(t) = T_{16}^0 + \int_0^t \left[ (b_{16})^{(2)} T_{17}(s_{(16)}) - \left( (b_{16}')^{(2)} - (b_{16}'')^{(2)} (G_{19}(s_{(16)}), s_{(16)}) \right) T_{16}(s_{(16)}) \right] ds_{(16)}$ | |
| $\bar{T}_{17}(t) = T_{17}^0 + \int_0^t \left[ (b_{17})^{(2)} T_{16}(s_{(16)}) - \left( (b_{17}')^{(2)} - (b_{17}'')^{(2)} (G_{19}(s_{(16)}), s_{(16)}) \right) T_{17}(s_{(16)}) \right] ds_{(16)}$ | |
| $\bar{T}_{18}(t) = T_{18}^0 + \int_0^t \left[ (b_{18})^{(2)} T_{17}(s_{(16)}) - \left( (b_{18}')^{(2)} - (b_{18}'')^{(2)} (G_{19}(s_{(16)}), s_{(16)}) \right) T_{18}(s_{(16)}) \right] ds_{(16)}$<br>Where $s_{(16)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(3)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{20})^{(3)}$, $T_i^0 \le (\hat{Q}_{20})^{(3)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{20})^{(3)} e^{(\bar{M}_{20})^{(3)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{20})^{(3)} e^{(\bar{M}_{20})^{(3)} t}$ | |
| By<br>$\bar{G}_{20}(t) = G_{20}^0 + \int_0^t \left[ (a_{20})^{(3)} G_{21}(s_{(20)}) - \left( (a_{20}')^{(3)} + a_{20}'')^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{20}(s_{(20)}) \right] ds_{(20)}$ | 161 |
| $\bar{G}_{21}(t) = G_{21}^0 + \int_0^t \left[ (a_{21})^{(3)} G_{20}(s_{(20)}) - \left( (a_{21}')^{(3)} + (a_{21}'')^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{21}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{G}_{22}(t) = G_{22}^0 + \int_0^t \left[ (a_{22})^{(3)} G_{21}(s_{(20)}) - \left( (a_{22}')^{(3)} + (a_{22}'')^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{22}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{20}(t) = T_{20}^0 + \int_0^t \left[ (b_{20})^{(3)} T_{21}(s_{(20)}) - \left( (b_{20}')^{(3)} - (b_{20}'')^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{20}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{21}(t) = T_{21}^0 + \int_0^t \left[ (b_{21})^{(3)} T_{20}(s_{(20)}) - \left( (b_{21}')^{(3)} - (b_{21}'')^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{21}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{22}(t) = T_{22}^0 + \int_0^t \left[ (b_{22})^{(3)} T_{21}(s_{(20)}) - \left( (b_{22}')^{(3)} - (b_{22}'')^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{22}(s_{(20)}) \right] ds_{(20)}$<br>Where $s_{(20)}$ is the integrand that is integrated over an interval $(0, t)$ | |

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| **Proof:** Consider operator $\mathcal{A}^{(4)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i$: $\mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{24})^{(4)}$ , $T_i^0 \leq (\hat{Q}_{24})^{(4)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ | |
| By<br>$\bar{G}_{24}(t) = G_{24}^0 + \int_0^t \left[ (a_{24})^{(4)} G_{25}(s_{(24)}) - \left( (a'_{24})^{(4)} + a''_{24})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{24}(s_{(24)}) \right] ds_{(24)}$ | 162 |
| $\bar{G}_{25}(t) = G_{25}^0 + \int_0^t \left[ (a_{25})^{(4)} G_{24}(s_{(24)}) - \left( (a'_{25})^{(4)} + (a''_{25})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{25}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{G}_{26}(t) = G_{26}^0 + \int_0^t \left[ (a_{26})^{(4)} G_{25}(s_{(24)}) - \left( (a'_{26})^{(4)} + (a''_{26})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{26}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{24}(t) = T_{24}^0 + \int_0^t \left[ (b_{24})^{(4)} T_{25}(s_{(24)}) - \left( (b'_{24})^{(4)} - (b''_{24})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{24}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{25}(t) = T_{25}^0 + \int_0^t \left[ (b_{25})^{(4)} T_{24}(s_{(24)}) - \left( (b'_{25})^{(4)} - (b''_{25})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{25}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{26}(t) = T_{26}^0 + \int_0^t \left[ (b_{26})^{(4)} T_{25}(s_{(24)}) - \left( (b'_{26})^{(4)} - (b''_{26})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{26}(s_{(24)}) \right] ds_{(24)}$<br>Where $s_{(24)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:** Consider operator $\mathcal{A}^{(5)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i$: $\mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{28})^{(5)}$ , $T_i^0 \leq (\hat{Q}_{28})^{(5)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| By<br>$\bar{G}_{28}(t) = G_{28}^0 + \int_0^t \left[ (a_{28})^{(5)} G_{29}(s_{(28)}) - \left( (a'_{28})^{(5)} + a''_{28})^{(5)} (T_{29}(s_{(28)}), s_{(28)}) \right) G_{28}(s_{(28)}) \right] ds_{(28)}$ | 163 |
| $\bar{G}_{29}(t) = G_{29}^0 + \int_0^t \left[ (a_{29})^{(5)} G_{28}(s_{(28)}) - \left( (a'_{29})^{(5)} + (a''_{29})^{(5)} (T_{29}(s_{(28)}), s_{(28)}) \right) G_{29}(s_{(28)}) \right] ds_{(28)}$ | |
| $\bar{G}_{30}(t) = G_{30}^0 + \int_0^t \left[ (a_{30})^{(5)} G_{29}(s_{(28)}) - \left( (a'_{30})^{(5)} + (a''_{30})^{(5)} (T_{29}(s_{(28)}), s_{(28)}) \right) G_{30}(s_{(28)}) \right] ds_{(28)}$ | |
| $\bar{T}_{28}(t) = T_{28}^0 + \int_0^t \left[ (b_{28})^{(5)} T_{29}(s_{(28)}) - \left( (b'_{28})^{(5)} - (b''_{28})^{(5)} (G_{31}(s_{(28)}), s_{(28)}) \right) T_{28}(s_{(28)}) \right] ds_{(28)}$ | |

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| $\bar{T}_{29}(t) = T_{29}^0 + \int_0^t \left[(b_{29})^{(5)} T_{28}(s_{(28)}) - \left((b'_{29})^{(5)} - (b''_{29})^{(5)}(G_{31}(s_{(28)}), s_{(28)})\right) T_{29}(s_{(28)})\right] ds_{(28)}$ | |
| $\bar{\mathrm{T}}_{30}(\mathrm{t}) = \mathrm{T}_{30}^0 + \int_0^t \left[(b_{30})^{(5)} T_{29}(s_{(28)}) - \left((b'_{30})^{(5)} - (b''_{30})^{(5)}(G_{31}(s_{(28)}), s_{(28)})\right) T_{30}(s_{(28)})\right] ds_{(28)}$<br>Where $s_{(28)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(6)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \le (\hat{P}_{32})^{(6)}$ , $T_i^0 \le (\hat{Q}_{32})^{(6)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ | |
| By<br>$\bar{G}_{32}(t) = G_{32}^0 + \int_0^t \left[(a_{32})^{(6)} G_{33}(s_{(32)}) - \left((a'_{32})^{(6)} + a''_{32})^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{32}(s_{(32)})\right] ds_{(32)}$ | 164 |
| $\bar{G}_{33}(t) = G_{33}^0 + \int_0^t \left[(a_{33})^{(6)} G_{32}(s_{(32)}) - \left((a'_{33})^{(6)} + (a''_{33})^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{33}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{G}_{34}(t) = G_{34}^0 + \int_0^t \left[(a_{34})^{(6)} G_{33}(s_{(32)}) - \left((a'_{34})^{(6)} + (a''_{34})^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{34}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{T}_{32}(t) = T_{32}^0 + \int_0^t \left[(b_{32})^{(6)} T_{33}(s_{(32)}) - \left((b'_{32})^{(6)} - (b''_{32})^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{32}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{T}_{33}(t) = T_{33}^0 + \int_0^t \left[(b_{33})^{(6)} T_{32}(s_{(32)}) - \left((b'_{33})^{(6)} - (b''_{33})^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{33}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{\mathrm{T}}_{34}(\mathrm{t}) = \mathrm{T}_{34}^0 + \int_0^t \left[(b_{34})^{(6)} T_{33}(s_{(32)}) - \left((b'_{34})^{(6)} - (b''_{34})^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{34}(s_{(32)})\right] ds_{(32)}$<br>Where $s_{(32)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(7)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \le (\hat{P}_{36})^{(7)}$ , $T_i^0 \le (\hat{Q}_{36})^{(7)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ | |

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| By<br>$\bar{G}_{36}(t) = G_{36}^{0} + \int_0^t \left[(a_{36})^{(7)} G_{37}(s_{(36)}) - \left((a_{36}^{'})^{(7)} + a_{36}^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{36}(s_{(36)})\right] ds_{(36)}$ | 165 |
| $\bar{G}_{37}(t) = G_{37}^{0} + \int_0^t \left[(a_{37})^{(7)} G_{36}(s_{(36)}) - \left((a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{37}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{G}_{38}(t) = G_{38}^{0} + \int_0^t \left[(a_{38})^{(7)} G_{37}(s_{(36)}) - \left((a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{38}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{36}(t) = T_{36}^{0} + \int_0^t \left[(b_{36})^{(7)} T_{37}(s_{(36)}) - \left((b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{36}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{37}(t) = T_{37}^{0} + \int_0^t \left[(b_{37})^{(7)} T_{36}(s_{(36)}) - \left((b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{37}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{38}(t) = T_{38}^{0} + \int_0^t \left[(b_{38})^{(7)} T_{37}(s_{(36)}) - \left((b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{38}(s_{(36)})\right] ds_{(36)}$<br>Where $s_{(36)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(8)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{40})^{(8)}$, $T_i^0 \le (\hat{Q}_{40})^{(8)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{40})^{(8)} e^{(\bar{M}_{40})^{(8)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{40})^{(8)} e^{(\bar{M}_{40})^{(8)} t}$ | |
| By<br>$\bar{G}_{40}(t) = G_{40}^{0} + \int_0^t \left[(a_{40})^{(8)} G_{41}(s_{(40)}) - \left((a_{40}^{'})^{(8)} + a_{40}^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\right) G_{40}(s_{(40)})\right] ds_{(40)}$ | 166 |
| $\bar{G}_{41}(t) = G_{41}^{0} + \int_0^t \left[(a_{41})^{(8)} G_{40}(s_{(40)}) - \left((a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\right) G_{41}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{G}_{42}(t) = G_{42}^{0} + \int_0^t \left[(a_{42})^{(8)} G_{41}(s_{(40)}) - \left((a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\right) G_{42}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{T}_{40}(t) = T_{40}^{0} + \int_0^t \left[(b_{40})^{(8)} T_{41}(s_{(40)}) - \left((b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}(G_{43}(s_{(40)}), s_{(40)})\right) T_{40}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{T}_{41}(t) = T_{41}^{0} + \int_0^t \left[(b_{41})^{(8)} T_{40}(s_{(40)}) - \left((b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}(G_{43}(s_{(40)}), s_{(40)})\right) T_{41}(s_{(40)})\right] ds_{(40)}$ | |

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| $\overline{T}_{42}(t) = T_{42}^0 + \int_0^t \left[(b_{42})^{(8)} T_{41}(s_{(40)}) - \left((b'_{42})^{(8)} - (b''_{42})^{(8)}(G_{43}(s_{(40)}), s_{(40)})\right) T_{42}(s_{(40)})\right] ds_{(40)}$<br><br>Where $s_{(40)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(9)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 166<br>A |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{44})^{(9)}$, $T_i^0 \le (\hat{Q}_{44})^{(9)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{44})^{(9)} e^{(\bar{M}_{44})^{(9)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{44})^{(9)} e^{(\bar{M}_{44})^{(9)} t}$ | |
| By<br><br>$\bar{G}_{44}(t) = G_{44}^0 + \int_0^t \left[(a_{44})^{(9)} G_{45}(s_{(44)}) - \left((a'_{44})^{(9)} + a''_{44})^{(9)}(T_{45}(s_{(44)}), s_{(44)})\right) G_{44}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{G}_{45}(t) = G_{45}^0 + \int_0^t \left[(a_{45})^{(9)} G_{44}(s_{(44)}) - \left((a'_{45})^{(9)} + (a''_{45})^{(9)}(T_{45}(s_{(44)}), s_{(44)})\right) G_{45}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{G}_{46}(t) = G_{46}^0 + \int_0^t \left[(a_{46})^{(9)} G_{45}(s_{(44)}) - \left((a'_{46})^{(9)} + (a''_{46})^{(9)}(T_{45}(s_{(44)}), s_{(44)})\right) G_{46}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{44}(t) = T_{44}^0 + \int_0^t \left[(b_{44})^{(9)} T_{45}(s_{(44)}) - \left((b'_{44})^{(9)} - (b''_{44})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{44}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{45}(t) = T_{45}^0 + \int_0^t \left[(b_{45})^{(9)} T_{44}(s_{(44)}) - \left((b'_{45})^{(9)} - (b''_{45})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{45}(s_{(44)})\right] ds_{(44)}$ | |
| $\overline{T}_{46}(t) = T_{46}^0 + \int_0^t \left[(b_{46})^{(9)} T_{45}(s_{(44)}) - \left((b'_{46})^{(9)} - (b''_{46})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{46}(s_{(44)})\right] ds_{(44)}$<br><br>Where $s_{(44)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| The operator $\mathcal{A}^{(1)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{13}(t) \le G_{13}^0 + \int_0^t \left[(a_{13})^{(1)} \left(G_{14}^0 + (\hat{P}_{13})^{(1)} e^{(\bar{M}_{13})^{(1)} s_{(13)}}\right)\right] ds_{(13)} =$<br><br>$\left(1 + (a_{13})^{(1)} t\right) G_{14}^0 + \frac{(a_{13})^{(1)} (\hat{P}_{13})^{(1)}}{(\bar{M}_{13})^{(1)}} \left(e^{(\bar{M}_{13})^{(1)} t} - 1\right)$ | 167 |
| From which it follows that<br><br>$(G_{13}(t) - G_{13}^0) e^{-(\bar{M}_{13})^{(1)} t} \le \frac{(a_{13})^{(1)}}{(\bar{M}_{13})^{(1)}} \left[\left((\hat{P}_{13})^{(1)} + G_{14}^0\right) e^{\left(-\frac{(\hat{P}_{13})^{(1)} + G_{14}^0}{G_{14}^0}\right)} + (\hat{P}_{13})^{(1)}\right]$<br><br>$(G_i^0)$ is as defined in the statement of theorem 1 | 168 |
| Analogous inequalities hold also for $G_{14}, G_{15}, T_{13}, T_{14}, T_{15}$ | |

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|---|---|
| The operator $\mathcal{A}^{(2)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that | |
| $G_{16}(t) \leq G_{16}^0 + \int_0^t \left[(a_{16})^{(2)} \left(G_{17}^0 + (\hat{P}_{16})^{(6)} e^{(\hat{M}_{16})^{(2)} s_{(16)}}\right)\right] ds_{(16)} =$ $\left(1 + (a_{16})^{(2)} t\right) G_{17}^0 + \frac{(a_{16})^{(2)} (\hat{P}_{16})^{(2)}}{(\hat{M}_{16})^{(2)}} \left(e^{(\hat{M}_{16})^{(2)} t} - 1\right)$ | 169 |
| From which it follows that $(G_{16}(t) - G_{16}^0) e^{-(\hat{M}_{16})^{(2)} t} \leq \frac{(a_{16})^{(2)}}{(\hat{M}_{16})^{(2)}} \left[\left((\hat{P}_{16})^{(2)} + G_{17}^0\right) e^{\left(-\frac{(\hat{P}_{16})^{(2)} + G_{17}^0}{G_{17}^0}\right)} + (\hat{P}_{16})^{(2)}\right]$ | 170 |
| Analogous inequalities hold also for $G_{17}, G_{18}, T_{16}, T_{17}, T_{18}$ | |
| The operator $\mathcal{A}^{(3)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{20}(t) \leq G_{20}^0 + \int_0^t \left[(a_{20})^{(3)} \left(G_{21}^0 + (\hat{P}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} s_{(20)}}\right)\right] ds_{(20)} =$ $\left(1 + (a_{20})^{(3)} t\right) G_{21}^0 + \frac{(a_{20})^{(3)} (\hat{P}_{20})^{(3)}}{(\hat{M}_{20})^{(3)}} \left(e^{(\hat{M}_{20})^{(3)} t} - 1\right)$ | 171 |
| From which it follows that $(G_{20}(t) - G_{20}^0) e^{-(\hat{M}_{20})^{(3)} t} \leq \frac{(a_{20})^{(3)}}{(\hat{M}_{20})^{(3)}} \left[\left((\hat{P}_{20})^{(3)} + G_{21}^0\right) e^{\left(-\frac{(\hat{P}_{20})^{(3)} + G_{21}^0}{G_{21}^0}\right)} + (\hat{P}_{20})^{(3)}\right]$ | 172 |
| Analogous inequalities hold also for $G_{21}, G_{22}, T_{20}, T_{21}, T_{22}$ | |
| The operator $\mathcal{A}^{(4)}$ maps the space of functions satisfying into itself .Indeed it is obvious that $G_{24}(t) \leq G_{24}^0 + \int_0^t \left[(a_{24})^{(4)} \left(G_{25}^0 + (\hat{P}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} s_{(24)}}\right)\right] ds_{(24)} =$ $\left(1 + (a_{24})^{(4)} t\right) G_{25}^0 + \frac{(a_{24})^{(4)} (\hat{P}_{24})^{(4)}}{(\hat{M}_{24})^{(4)}} \left(e^{(\hat{M}_{24})^{(4)} t} - 1\right)$ | 173 |
| From which it follows that $(G_{24}(t) - G_{24}^0) e^{-(\hat{M}_{24})^{(4)} t} \leq \frac{(a_{24})^{(4)}}{(\hat{M}_{24})^{(4)}} \left[\left((\hat{P}_{24})^{(4)} + G_{25}^0\right) e^{\left(-\frac{(\hat{P}_{24})^{(4)} + G_{25}^0}{G_{25}^0}\right)} + (\hat{P}_{24})^{(4)}\right]$ $(G_i^0)$ is as defined in the statement of theorem 4 | 174 |
| The operator $\mathcal{A}^{(5)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{28}(t) \leq G_{28}^0 + \int_0^t \left[(a_{28})^{(5)} \left(G_{29}^0 + (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} s_{(28)}}\right)\right] ds_{(28)} =$ $\left(1 + (a_{28})^{(5)} t\right) G_{29}^0 + \frac{(a_{28})^{(5)} (\hat{P}_{28})^{(5)}}{(\hat{M}_{28})^{(5)}} \left(e^{(\hat{M}_{28})^{(5)} t} - 1\right)$ | |

From which it follows that

$$(G_{28}(t)-G_{28}^{0})e^{-(\bar{M}_{28})^{(5)}t}\leq\frac{(a_{28})^{(5)}}{(\bar{M}_{28})^{(5)}}\left[((\hat{P}_{28})^{(5)}+G_{29}^{0})e^{\left(-\frac{(\hat{P}_{28})^{(5)}+G_{29}^{0}}{G_{29}^{0}}\right)}+(\hat{P}_{28})^{(5)}\right]$$

$(G_i^0)$ is as defined in the statement of theorem 5 175

The operator $\mathcal{A}^{(6)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that 176

$$G_{32}(t)\leq G_{32}^{0}+\int_{0}^{t}\left[(a_{32})^{(6)}\left(G_{33}^{0}+(\hat{P}_{32})^{(6)}e^{(\bar{M}_{32})^{(6)}s_{(32)}}\right)\right]ds_{(32)}=$$

$$\left(1+(a_{32})^{(6)}t\right)G_{33}^{0}+\frac{(a_{32})^{(6)}(\hat{P}_{32})^{(6)}}{(\bar{M}_{32})^{(6)}}\left(e^{(\bar{M}_{32})^{(6)}t}-1\right)$$

From which it follows that 177

$$(G_{32}(t)-G_{32}^{0})e^{-(\bar{M}_{32})^{(6)}t}\leq\frac{(a_{32})^{(6)}}{(\bar{M}_{32})^{(6)}}\left[((\hat{P}_{32})^{(6)}+G_{33}^{0})e^{\left(-\frac{(\hat{P}_{32})^{(6)}+G_{33}^{0}}{G_{33}^{0}}\right)}+(\hat{P}_{32})^{(6)}\right]$$

$(G_i^0)$ is as defined in the statement of theorem 6

Analogous inequalities hold also for $G_{25}, G_{26}, T_{24}, T_{25}, T_{26}$

(f) The operator $\mathcal{A}^{(7)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that 178

$$G_{36}(t)\leq G_{36}^{0}+\int_{0}^{t}\left[(a_{36})^{(7)}\left(G_{37}^{0}+(\hat{P}_{36})^{(7)}e^{(\bar{M}_{36})^{(7)}s_{(36)}}\right)\right]ds_{(36)}=$$

$$\left(1+(a_{36})^{(7)}t\right)G_{37}^{0}+\frac{(a_{36})^{(7)}(\hat{P}_{36})^{(7)}}{(\bar{M}_{36})^{(7)}}\left(e^{(\bar{M}_{36})^{(7)}t}-1\right)$$

From which it follows that

$$(G_{36}(t)-G_{36}^{0})e^{-(\bar{M}_{36})^{(7)}t}\leq\frac{(a_{36})^{(7)}}{(\bar{M}_{36})^{(7)}}\left[((\hat{P}_{36})^{(7)}+G_{37}^{0})e^{\left(-\frac{(\hat{P}_{36})^{(7)}+G_{37}^{0}}{G_{37}^{0}}\right)}+(\hat{P}_{36})^{(7)}\right]$$

$(G_i^0)$ is as defined in the statement of theorem 7

The operator $\mathcal{A}^{(8)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that

$$G_{40}(t)\leq G_{40}^{0}+\int_{0}^{t}\left[(a_{40})^{(8)}\left(G_{41}^{0}+(\hat{P}_{40})^{(8)}e^{(\bar{M}_{40})^{(8)}s_{(40)}}\right)\right]ds_{(40)}=$$

$$\left(1+(a_{40})^{(8)}t\right)G_{41}^{0}+\frac{(a_{40})^{(8)}(\hat{P}_{40})^{(8)}}{(\bar{M}_{40})^{(8)}}\left(e^{(\bar{M}_{40})^{(8)}t}-1\right)$$ 180

From which it follows that

$$(G_{40}(t)-G_{40}^{0})e^{-(\bar{M}_{40})^{(8)}t}\leq\frac{(a_{40})^{(8)}}{(\bar{M}_{40})^{(8)}}\left[((\hat{P}_{40})^{(8)}+G_{41}^{0})e^{\left(-\frac{(\hat{P}_{40})^{(8)}+G_{41}^{0}}{G_{41}^{0}}\right)}+(\hat{P}_{40})^{(8)}\right]$$ 181

$(G_i^0)$ is as defined in the statement of theorem 8
Analogous inequalities hold also for $G_{41}, G_{42}, T_{40}, T_{41}, T_{42}$

The operator $\mathcal{A}^{(9)}$ maps the space of functions satisfying 34,35,36 into itself .Indeed it is obvious that

$$G_{44}(t) \le G_{44}^0 + \int_0^t \left[(a_{44})^{(9)}\left(G_{45}^0 + (\hat{P}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} s_{(44)}}\right)\right] ds_{(44)} =$$

$$\left(1 + (a_{44})^{(9)} t\right) G_{45}^0 + \frac{(a_{44})^{(9)} (\hat{P}_{44})^{(9)}}{(\hat{M}_{44})^{(9)}}\left(e^{(\hat{M}_{44})^{(9)} t} - 1\right)$$

From which it follows that

$$(G_{44}(t) - G_{44}^0) e^{-(\hat{M}_{44})^{(9)} t} \le \frac{(a_{44})^{(9)}}{(\hat{M}_{44})^{(9)}}\left[\left((\hat{P}_{44})^{(9)} + G_{45}^0\right) e^{\left(-\frac{(\hat{P}_{44})^{(9)} + G_{45}^0}{G_{45}^0}\right)} + (\hat{P}_{44})^{(9)}\right]$$

$(G_i^0)$ is as defined in the statement of theorem 9

Analogous inequalities hold also for $G_{45}, G_{46}, T_{44}, T_{45}, T_{46}$

It is now sufficient to take $\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}} < 1$ and to choose

$(\hat{P}_{13})^{(1)}$ and $(\hat{Q}_{13})^{(1)}$ large to have

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$$\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}\left[(\hat{P}_{13})^{(1)} + \left((\hat{P}_{13})^{(1)} + G_j^0\right) e^{-\left(\frac{(\hat{P}_{13})^{(1)} + G_j^0}{G_j^0}\right)}\right] \le (\hat{P}_{13})^{(1)} \quad (183)$$

$$\frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}}\left[\left((\hat{Q}_{13})^{(1)} + T_j^0\right) e^{-\left(\frac{(\hat{Q}_{13})^{(1)} + T_j^0}{T_j^0}\right)} + (\hat{Q}_{13})^{(1)}\right] \le (\hat{Q}_{13})^{(1)} \quad (184)$$

In order that the operator $\mathcal{A}^{(1)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself

The operator $\mathcal{A}^{(1)}$ is a contraction with respect to the metric (185)

$$d\left(\left(G^{(1)}, T^{(1)}\right), \left(G^{(2)}, T^{(2)}\right)\right) =$$

$$\sup_i \{\max_{t \in \mathbb{R}_+} \left|G_i^{(1)}(t) - G_i^{(2)}(t)\right| e^{-(\hat{M}_{13})^{(1)} t}, \max_{t \in \mathbb{R}_+} \left|T_i^{(1)}(t) - T_i^{(2)}(t)\right| e^{-(\hat{M}_{13})^{(1)} t}\}$$

Indeed if we denote

**Definition of** $\tilde{G}, \tilde{T}$ : $\left(\tilde{G}, \tilde{T}\right) = \mathcal{A}^{(1)}(G, T)$

It results

$$\left|\tilde{G}_{13}^{(1)}-\tilde{G}_{i}^{(2)}\right| \leq \int_0^t (a_{13})^{(1)} \left|G_{14}^{(1)}-G_{14}^{(2)}\right| e^{-(\widehat{M}_{13})^{(1)}s_{(13)}} e^{(\widehat{M}_{13})^{(1)}s_{(13)}}\, ds_{(13)} +$$

$$\int_0^t \{(a'_{13})^{(1)} \left|G_{13}^{(1)}-G_{13}^{(2)}\right| e^{-(\widehat{M}_{13})^{(1)}s_{(13)}} e^{-(\widehat{M}_{13})^{(1)}s_{(13)}} +$$

$$(a''_{13})^{(1)}\left(T_{14}^{(1)}, s_{(13)}\right)\left|G_{13}^{(1)}-G_{13}^{(2)}\right| e^{-(\widehat{M}_{13})^{(1)}s_{(13)}} e^{(\widehat{M}_{13})^{(1)}s_{(13)}} +$$

$$G_{13}^{(2)}\left|(a''_{13})^{(1)}\left(T_{14}^{(1)}, s_{(13)}\right)-(a''_{13})^{(1)}\left(T_{14}^{(2)}, s_{(13)}\right)\right|\ e^{-(\widehat{M}_{13})^{(1)}s_{(13)}} e^{(\widehat{M}_{13})^{(1)}s_{(13)}}\} ds_{(13)}$$

Where $s_{(13)}$ represents integrand that is integrated over the interval $[0, t]$

From the hypotheses it follows

$$\left|G^{(1)}-G^{(2)}\right| e^{-(\widehat{M}_{13})^{(1)}t} \leq$$
$$\frac{1}{(\widehat{M}_{13})^{(1)}}\left((a_{13})^{(1)} + (a'_{13})^{(1)} + (\widehat{A}_{13})^{(1)} + (\widehat{P}_{13})^{(1)}(\widehat{k}_{13})^{(1)}\right) d\left(\left(G^{(1)}, T^{(1)};\ G^{(2)}, T^{(2)}\right)\right)$$

And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows

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**Remark 1:** The fact that we supposed $(a''_{13})^{(1)}$ and $(b''_{13})^{(1)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{13})^{(1)} e^{(\widehat{M}_{13})^{(1)}t}$ *and* $(\widehat{Q}_{13})^{(1)} e^{(\widehat{M}_{13})^{(1)}t}$ respectively of $\mathbb{R}_+$.

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a''_i)^{(1)}$ and $(b''_i)^{(1)}, i = 13,14,15$ depend only on $T_{14}$ and respectively on $G(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition.

**Remark 2:** There does not exist any $t$ where $G_i(t) = 0$ *and* $T_i(t) = 0$

From 19 to 24 it results

$$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a'_i)^{(1)} - (a''_i)^{(1)}(T_{14}(s_{(13)}), s_{(13)})\} ds_{(13)}\right]} \geq 0$$

$$T_i(t) \geq T_i^0 e^{\left(-(b'_i)^{(1)}t\right)} > 0 \quad \text{for t} > 0$$

**Definition of** $\left((\widehat{M}_{13})^{(1)}\right)_1, \left((\widehat{M}_{13})^{(1)}\right)_2$ *and* $\left((\widehat{M}_{13})^{(1)}\right)_3$ :

**Remark 3**: if $G_{13}$ is bounded, the same property have also $G_{14}$ *and* $G_{15}$ . indeed if

$G_{13} < (\widehat{M}_{13})^{(1)}$ it follows $\frac{dG_{14}}{dt} \leq \left((\widehat{M}_{13})^{(1)}\right)_1 - (a'_{14})^{(1)} G_{14}$ and by integrating

$$G_{14} \leq \left((\widehat{M}_{13})^{(1)}\right)_2 = G_{14}^0 + 2(a_{14})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_1 / (a'_{14})^{(1)}$$

In the same way , one can obtain

$$G_{15} \leq \left((\widehat{M}_{13})^{(1)}\right)_3 = G_{15}^0 + 2(a_{15})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_2 / (a'_{15})^{(1)}$$

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| If $G_{14}$ *or* $G_{15}$ is bounded, the same property follows for $G_{13}$ , $G_{15}$ and $G_{13}$ , $G_{14}$ respectively. | |
| **Remark 4:** If $G_{13}$ *is* bounded, from below, the same property holds for $G_{14}$ *and* $G_{15}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{14}$ is bounded from below. | 188 |
| **Remark 5:** If $T_{13}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(1)}(G(t),t)) = (b_{14}^{'})^{(1)}$ then $T_{14} \to \infty$.<br><br>**Definition of** $(m)^{(1)}$ and $\varepsilon_1$ :<br><br>Indeed let $t_1$ be so that for $t > t_1$<br><br>$(b_{14})^{(1)} - (b_i^{''})^{(1)}(G(t),t) < \varepsilon_1, T_{13}(t) > (m)^{(1)}$ | 189 |
| Then $\frac{dT_{14}}{dt} \geq (a_{14})^{(1)}(m)^{(1)} - \varepsilon_1 T_{14}$ which leads to<br><br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{\varepsilon_1}\right)(1-e^{-\varepsilon_1 t}) + T_{14}^0 e^{-\varepsilon_1 t}$ If we take t such that $e^{-\varepsilon_1 t} = \frac{1}{2}$ it results<br><br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{2}\right)$, $t = log\frac{2}{\varepsilon_1}$ By taking now $\varepsilon_1$ sufficiently small one sees that $T_{14}$ is unbounded.<br>The same property holds for $T_{15}$ if $\lim_{t\to\infty}(b_{15}^{''})^{(1)}(G(t),t) = (b_{15}^{'})^{(1)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |
| It is now sufficient to take $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}$ , $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ and to choose<br><br>$(\hat{P}_{16})^{(2)}$ *and* $(\hat{Q}_{16})^{(2)}$ large to have | 190 |
| $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[(\hat{P}_{16})^{(2)} + \left((\hat{P}_{16})^{(2)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{16})^{(2)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{16})^{(2)}$ | 191 |
| $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[\left((\hat{Q}_{16})^{(2)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{16})^{(2)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{16})^{(2)}\right] \leq (\hat{Q}_{16})^{(2)}$ | 192 |
| In order that the operator $\mathcal{A}^{(2)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 193 |
| The operator $\mathcal{A}^{(2)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{19})^{(1)},(T_{19})^{(1)}),((G_{19})^{(2)},(T_{19})^{(2)})\right) =$<br><br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t) - G_i^{(2)}(t)\right|e^{-(\hat{M}_{16})^{(2)}t}, \max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t) - T_i^{(2)}(t)\right|e^{-(\hat{M}_{16})^{(2)}t}\}$ | 194 |
| Indeed if we denote | 195 |

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| **Definition of** $\widetilde{G_{19}}, \widetilde{T_{19}}$ : $\left(\widetilde{G_{19}}, \widetilde{T_{19}}\right) = \mathcal{A}^{(2)}(G_{19}, T_{19})$ | |
| It results<br>$\left\|\tilde{G}_{16}^{(1)} - \tilde{G}_{i}^{(2)}\right\| \leq \int_0^t (a_{16})^{(2)} \left\|G_{17}^{(1)} - G_{17}^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{(\widehat{M}_{16})^{(2)} s_{(16)}}\, ds_{(16)} +$<br>$\int_0^t \{(a_{16}^{'})^{(2)} \left\|G_{16}^{(1)} - G_{16}^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} +$<br>$(a_{16}^{''})^{(2)}\left(T_{17}^{(1)}, s_{(16)}\right)\left\|G_{16}^{(1)} - G_{16}^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{(\widehat{M}_{16})^{(2)} s_{(16)}} +$<br>$G_{16}^{(2)} \left\|(a_{16}^{''})^{(2)}\left(T_{17}^{(1)}, s_{(16)}\right) - (a_{16}^{''})^{(2)}\left(T_{17}^{(2)}, s_{(16)}\right)\right\| \; e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{(\widehat{M}_{16})^{(2)} s_{(16)}} \} ds_{(16)}$ | 196 |
| Where $s_{(16)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | 197 |
| $\left\|(G_{19})^{(1)} - (G_{19})^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} t} \leq$<br>$\frac{1}{(\widehat{M}_{16})^{(2)}}\left((a_{16})^{(2)} + (a_{16}^{'})^{(2)} + (\widehat{A}_{16})^{(2)} + (\widehat{P}_{16})^{(2)}(\widehat{k}_{16})^{(2)}\right) d\left(((G_{19})^{(1)}, (T_{19})^{(1)}; (G_{19})^{(2)}, (T_{19})^{(2)})\right)$ | |
| And analogous inequalities for $G_i$ and $T_i$. Taking into account the hypothesis the result follows | 198 |
| **Remark 6:** The fact that we supposed $(a_{16}^{''})^{(2)}$ and $(b_{16}^{''})^{(2)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{16})^{(2)} e^{(\widehat{M}_{16})^{(2)} t}$ and $(\widehat{Q}_{16})^{(2)} e^{(\widehat{M}_{16})^{(2)} t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}, i = 16,17,18$ depend only on $T_{17}$ and respectively on $(G_{19})$(and not on t) and hypothesis can replaced by a usual Lipschitz condition. | 199 |
| **Remark 7:** There does not exist any t where $G_i(t) = 0$ and $T_i(t) = 0$<br>it results<br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(2)} - (a_i^{''})^{(2)}(T_{17}(s_{(16)}), s_{(16)})\} ds_{(16)}\right]} \geq 0$<br>$T_i(t) \geq T_i^0 e^{(-(b_i^{'})^{(2)} t)} > 0$ for $t > 0$ | 200 |
| **Definition of** $\left((\widehat{M}_{16})^{(2)}\right)_1$, $\left((\widehat{M}_{16})^{(2)}\right)_2$ and $\left((\widehat{M}_{16})^{(2)}\right)_3$ :<br>**Remark 8:** if $G_{16}$ is bounded, the same property have also $G_{17}$ and $G_{18}$ . indeed if<br>$G_{16} < (\widehat{M}_{16})^{(2)}$ it follows $\frac{dG_{17}}{dt} \leq \left((\widehat{M}_{16})^{(2)}\right)_1 - (a_{17}^{'})^{(2)} G_{17}$ and by integrating<br>$G_{17} \leq \left((\widehat{M}_{16})^{(2)}\right)_2 = G_{17}^0 + 2(a_{17})^{(2)}\left((\widehat{M}_{16})^{(2)}\right)_1 / (a_{17}^{'})^{(2)}$ | 201 |

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| In the same way , one can obtain<br>$G_{18} \leq \left((\widehat{M}_{16})^{(2)}\right)_3 = G_{18}^0 + 2(a_{18})^{(2)}\left((\widehat{M}_{16})^{(2)}\right)_2/(a_{18}')^{(2)}$<br>If $G_{17}$ or $G_{18}$ is bounded, the same property follows for $G_{16}$ , $G_{18}$ and $G_{16}$ , $G_{17}$ respectively. | |
| **Remark 9:** If $G_{16}$ is bounded, from below, the same property holds for $G_{17}$ and $G_{18}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{17}$ is bounded from below. | 202 |
| **Remark 10:** If $T_{16}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(2)}((G_{19})(t),t)) = (b_{17}')^{(2)}$ then $T_{17} \to \infty$.<br>**Definition of** $(m)^{(2)}$ and $\varepsilon_2$ :<br>Indeed let $t_2$ be so that for $t > t_2$<br>$(b_{17})^{(2)} - (b_i'')^{(2)}((G_{19})(t),t) < \varepsilon_2, T_{16}(t) > (m)^{(2)}$ | 203 |
| Then $\frac{dT_{17}}{dt} \geq (a_{17})^{(2)}(m)^{(2)} - \varepsilon_2 T_{17}$ which leads to<br>$T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{\varepsilon_2}\right)(1 - e^{-\varepsilon_2 t}) + T_{17}^0 e^{-\varepsilon_2 t}$ If we take t such that $e^{-\varepsilon_2 t} = \frac{1}{2}$ it results | 204 |
| $T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{2}\right)$, $t = \log\frac{2}{\varepsilon_2}$ By taking now $\varepsilon_2$ sufficiently small one sees that $T_{17}$ is unbounded.<br>The same property holds for $T_{18}$ if $\lim_{t\to\infty}(b_{18}'')^{(2)}\left((G_{19})(t),t\right) = (b_{18}')^{(2)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 205 |
| It is now sufficient to take $\frac{(a_i)^{(3)}}{(\widehat{M}_{20})^{(3)}}$ , $\frac{(b_i)^{(3)}}{(\widehat{M}_{20})^{(3)}} < 1$ and to choose<br>$(\widehat{P}_{20})^{(3)}$ and $(\widehat{Q}_{20})^{(3)}$ large to have | 207 |
| $\frac{(a_i)^{(3)}}{(\widehat{M}_{20})^{(3)}}\left[(\widehat{P}_{20})^{(3)} + \left((\widehat{P}_{20})^{(3)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{20})^{(3)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{20})^{(3)}$ | 208 |
| $\frac{(b_i)^{(3)}}{(\widehat{M}_{20})^{(3)}}\left[\left((\widehat{Q}_{20})^{(3)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{20})^{(3)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{20})^{(3)}\right] \leq (\widehat{Q}_{20})^{(3)}$ | 209 |
| In order that the operator $\mathcal{A}^{(3)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 210 |
| The operator $\mathcal{A}^{(3)}$ is a contraction with respect to the metric<br>$d\left(((G_{23})^{(1)},(T_{23})^{(1)}),((G_{23})^{(2)},(T_{23})^{(2)})\right) =$ | 211 |

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| $\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{20})^{(3)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{20})^{(3)}t}\}$ | |
| Indeed if we denote<br>**Definition of** $\widetilde{G_{23}},\widetilde{T_{23}}:(\widetilde{(G_{23})},\widetilde{(T_{23})})=\mathcal{A}^{(3)}((G_{23}),(T_{23}))$ | 212 |
| It results<br>$\left\|\tilde{G}_{20}^{(1)}-\tilde{G}_i^{(2)}\right\|\le\int_0^t(a_{20})^{(3)}\left\|G_{21}^{(1)}-G_{21}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{(\widehat{M}_{20})^{(3)}s_{(20)}}\,ds_{(20)}+$<br>$\int_0^t\{(a_{20}^{'})^{(3)}\left\|G_{20}^{(1)}-G_{20}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}+$<br>$(a_{20}^{''})^{(3)}\left(T_{21}^{(1)},s_{(20)}\right)\left\|G_{20}^{(1)}-G_{20}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{(\widehat{M}_{20})^{(3)}s_{(20)}}+$<br>$G_{20}^{(2)}\left\|(a_{20}^{''})^{(3)}\left(T_{21}^{(1)},s_{(20)}\right)-(a_{20}^{''})^{(3)}\left(T_{21}^{(2)},s_{(20)}\right)\right\|\ e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{(\widehat{M}_{20})^{(3)}s_{(20)}}\}ds_{(20)}$<br>Where $s_{(20)}$ represents integrand that is integrated over the interval $[0,t]$<br>From the hypotheses it follows | 213 |
| $\left\|G_{23}{}^{(1)}-G_{23}{}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}t}\le$<br>$\frac{1}{(\widehat{M}_{20})^{(3)}}\left((a_{20})^{(3)}+(a_{20}^{'})^{(3)}+(\widehat{A}_{20})^{(3)}+(\widehat{P}_{20})^{(3)}(\widehat{k}_{20})^{(3)}\right)d\left(((G_{23})^{(1)},(T_{23})^{(1)};\ (G_{23})^{(2)},(T_{23})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 214 |
| **Remark 11:** The fact that we supposed $(a_{20}^{''})^{(3)}$ and $(b_{20}^{''})^{(3)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}$ *and* $(\widehat{Q}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(3)}$ and $(b_i^{''})^{(3)}, i=20,21,22$ depend only on $T_{21}$ and respectively on $(G_{23})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 215 |
| **Remark 12:** There does not exist any $t$ where $G_i(t)=0$ *and* $T_i(t)=0$<br>it results<br>$G_i(t)\ge G_i^0e^{\left[-\int_0^t\{(a_i^{'})^{(3)}-(a_i^{''})^{(3)}(T_{21}(s_{(20)}),s_{(20)})\}ds_{(20)}\right]}\ge 0$<br>$T_i(t)\ge T_i^0e^{(-(b_i^{'})^{(3)}t)}>0$ for t > 0 | 216 |

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| **Definition of** $\left((\widehat{M}_{20})^{(3)}\right)_1, \left((\widehat{M}_{20})^{(3)}\right)_2$ *and* $\left((\widehat{M}_{20})^{(3)}\right)_3$ : | 217 |
| **Remark 13:** if $G_{20}$ is bounded, the same property have also $G_{21}$ *and* $G_{22}$ . indeed if<br>$G_{20} < (\widehat{M}_{20})^{(3)}$ it follows $\frac{dG_{21}}{dt} \leq \left((\widehat{M}_{20})^{(3)}\right)_1 - (a_{21}^{'})^{(3)} G_{21}$ and by integrating<br>$G_{21} \leq \left((\widehat{M}_{20})^{(3)}\right)_2 = G_{21}^0 + 2(a_{21})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_1/(a_{21}^{'})^{(3)}$<br>In the same way , one can obtain<br>$G_{22} \leq \left((\widehat{M}_{20})^{(3)}\right)_3 = G_{22}^0 + 2(a_{22})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_2/(a_{22}^{'})^{(3)}$<br>If $G_{21}$ *or* $G_{22}$ is bounded, the same property follows for $G_{20}$ , $G_{22}$ and $G_{20}$ , $G_{21}$ respectively. | |
| **Remark 14:** If $G_{20}$ *is* bounded, from below, the same property holds for $G_{21}$ *and* $G_{22}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{21}$ is bounded from below. | 218 |
| **Remark 15:** If $T_{20}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(3)}\left((G_{23})(t),t\right) = (b_{21}^{'})^{(3)}$ then $T_{21} \to \infty$.<br>**Definition of** $(m)^{(3)}$ and $\varepsilon_3$ :<br>Indeed let $t_3$ be so that for $t > t_3$<br>$(b_{21})^{(3)} - (b_i^{''})^{(3)}\left((G_{23})(t),t\right) < \varepsilon_3, T_{20}(t) > (m)^{(3)}$ | 219 |
| Then $\frac{dT_{21}}{dt} \geq (a_{21})^{(3)}(m)^{(3)} - \varepsilon_3 T_{21}$ which leads to<br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{\varepsilon_3}\right)(1 - e^{-\varepsilon_3 t}) + T_{21}^0 e^{-\varepsilon_3 t}$ If we take t such that $e^{-\varepsilon_3 t} = \frac{1}{2}$ it results<br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{2}\right)$, $t = log\frac{2}{\varepsilon_3}$ By taking now $\varepsilon_3$ sufficiently small one sees that $T_{21}$ is unbounded.<br>The same property holds for $T_{22}$ if $\lim_{t\to\infty}(b_{22}^{''})^{(3)}\left((G_{23})(t),t\right) = (b_{22}^{'})^{(3)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 220 |
| It is now sufficient to take $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}$ , $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}} < 1$ and to choose<br>$(\widehat{P}_{24})^{(4)}$ and $(\widehat{Q}_{24})^{(4)}$ large to have | 221 |
| $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[(\widehat{P}_{24})^{(4)} + \left((\widehat{P}_{24})^{(4)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{24})^{(4)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{24})^{(4)}$ | 222 |
| $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[\left((\widehat{Q}_{24})^{(4)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{24})^{(4)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{24})^{(4)}\right] \leq (\widehat{Q}_{24})^{(4)}$ | 223 |

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| In order that the operator $\mathcal{A}^{(4)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 224 |
| The operator $\mathcal{A}^{(4)}$ is a contraction with respect to the metric<br>$d\left(((G_{27})^{(1)},(T_{27})^{(1)}),((G_{27})^{(2)},(T_{27})^{(2)})\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{24})^{(4)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{24})^{(4)}t}\}$<br>Indeed if we denote<br>**Definition of** $(\widetilde{G_{27}}),(\widetilde{T_{27}})$ : $\left((\widetilde{G_{27}}),(\widetilde{T_{27}})\right)=\mathcal{A}^{(4)}((G_{27}),(T_{27}))$<br>It results<br>$\left\|\tilde{G}_{24}^{(1)}-\tilde{G}_i^{(2)}\right\|\le\int_0^t(a_{24})^{(4)}\left\|G_{25}^{(1)}-G_{25}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}\,ds_{(24)}+$<br>$\int_0^t\{(a_{24}^{'})^{(4)}\left\|G_{24}^{(1)}-G_{24}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}+$<br>$(a_{24}^{''})^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)\left\|G_{24}^{(1)}-G_{24}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}+$<br>$G_{24}^{(2)}\left\|(a_{24}^{''})^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)-(a_{24}^{''})^{(4)}\left(T_{25}^{(2)},s_{(24)}\right)\right\|\ e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}\}ds_{(24)}$<br>Where $s_{(24)}$ represents integrand that is integrated over the interval $[0,\mathrm{t}]$<br>From the hypotheses on Equations it follows | 225 |
| $\left\|(G_{27})^{(1)}-(G_{27})^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}t}\le$<br>$\frac{1}{(\widehat{M}_{24})^{(4)}}\left((a_{24})^{(4)}+(a_{24}^{'})^{(4)}+(\widehat{A}_{24})^{(4)}+(\widehat{P}_{24})^{(4)}(\widehat{k}_{24})^{(4)}\right)d\left(((G_{27})^{(1)},(T_{27})^{(1)};\ (G_{27})^{(2)},(T_{27})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 226 |
| **Remark 16:** The fact that we supposed $(a_{24}^{''})^{(4)}$ and $(b_{24}^{''})^{(4)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}$ *and* $(\widehat{Q}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(4)}$ and $(b_i^{''})^{(4)}, i=24{,}25{,}26$ depend only on $\mathrm{T}_{25}$ and respectively on $(G_{27})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 227 |
| **Remark 17:** There does not exist any $t$ where $G_i\,(t)=0\ and\ T_i\,(t)=0$<br>it results | 228 |

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| $G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i')^{(4)} - (a_i'')^{(4)}(T_{25}(s_{(24)}), s_{(24)})\} ds_{(24)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{\left(-(b_i')^{(4)} t\right)} > 0$ for t > 0 | |
| **Definition of** $\left((\widehat{M}_{24})^{(4)}\right)_1, \left((\widehat{M}_{24})^{(4)}\right)_2$ *and* $\left((\widehat{M}_{24})^{(4)}\right)_3$ :<br><br>**Remark 18:** if $G_{24}$ is bounded, the same property have also $G_{25}$ *and* $G_{26}$ . indeed if<br><br>$G_{24} < (\widehat{M}_{24})^{(4)}$ it follows $\frac{dG_{25}}{dt} \leq \left((\widehat{M}_{24})^{(4)}\right)_1 - (a_{25}')^{(4)} G_{25}$ and by integrating<br><br>$G_{25} \leq \left((\widehat{M}_{24})^{(4)}\right)_2 = G_{25}^0 + 2(a_{25})^{(4)} \left((\widehat{M}_{24})^{(4)}\right)_1 / (a_{25}')^{(4)}$<br><br>In the same way , one can obtain<br><br>$G_{26} \leq \left((\widehat{M}_{24})^{(4)}\right)_3 = G_{26}^0 + 2(a_{26})^{(4)} \left((\widehat{M}_{24})^{(4)}\right)_2 / (a_{26}')^{(4)}$<br><br>If $G_{25}$ *or* $G_{26}$ is bounded, the same property follows for $G_{24}$ , $G_{26}$ and $G_{24}$ , $G_{25}$ respectively. | 229 |
| **Remark 19:** If $G_{24}$ *is* bounded, from below, the same property holds for $G_{25}$ *and* $G_{26}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{25}$ is bounded from below. | 230 |
| **Remark 20:** If $T_{24}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(4)}((G_{27})(t), t)) = (b_{25}')^{(4)}$ then $T_{25} \to \infty$.<br><br>**Definition of** $(m)^{(4)}$ and $\varepsilon_4$ :<br><br>Indeed let $t_4$ be so that for $t > t_4$<br><br>$(b_{25})^{(4)} - (b_i'')^{(4)}((G_{27})(t), t) < \varepsilon_4, T_{24}(t) > (m)^{(4)}$ | 231 |
| Then $\frac{dT_{25}}{dt} \geq (a_{25})^{(4)}(m)^{(4)} - \varepsilon_4 T_{25}$ which leads to<br><br>$T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{\varepsilon_4}\right)(1 - e^{-\varepsilon_4 t}) + T_{25}^0 e^{-\varepsilon_4 t}$ If we take t such that $e^{-\varepsilon_4 t} = \frac{1}{2}$ it results<br><br>$T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{2}\right)$, $t = log\frac{2}{\varepsilon_4}$ By taking now $\varepsilon_4$ sufficiently small one sees that $T_{25}$ is unbounded. The same property holds for $T_{26}$ if $\lim_{t\to\infty}(b_{26}'')^{(4)}\left((G_{27})(t), t\right) = (b_{26}')^{(4)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 42<br><br>Analogous inequalities hold also for $G_{29}$ , $G_{30}, T_{28}, T_{29}, T_{30}$ | 232 |
| It is now sufficient to take $\frac{(a_i)^{(5)}}{(\widehat{M}_{28})^{(5)}}$ , $\frac{(b_i)^{(5)}}{(\widehat{M}_{28})^{(5)}} < 1$ and to choose<br><br>$(\widehat{P}_{28})^{(5)}$ and $(\widehat{Q}_{28})^{(5)}$ large to have | 233 |

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| $\frac{(a_i)^{(5)}}{(\widehat{M}_{28})^{(5)}}\left[(\widehat{P}_{28})^{(5)}+\left((\hat{P}_{28})^{(5)}+G_j^0\right)e^{-\left(\frac{(\hat{P}_{28})^{(5)}+G_j^0}{G_j^0}\right)}\right]\leq(\hat{P}_{28})^{(5)}$ | 234 |
| $\frac{(b_i)^{(5)}}{(\widehat{M}_{28})^{(5)}}\left[\left((\hat{Q}_{28})^{(5)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{28})^{(5)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{28})^{(5)}\right]\leq(\hat{Q}_{28})^{(5)}$ | 235 |
| In order that the operator $\mathcal{A}^{(5)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(5)}$ is a contraction with respect to the metric<br>$d\left(((G_{31})^{(1)},(T_{31})^{(1)}),((G_{31})^{(2)},(T_{31})^{(2)})\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{28})^{(5)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{28})^{(5)}t}\}$<br>Indeed if we denote<br>**Definition of** $\widetilde{(G_{31})},\widetilde{(T_{31})}$ : $\left(\widetilde{(G_{31})},\widetilde{(T_{31})}\right)=\mathcal{A}^{(5)}\left((G_{31}),(T_{31})\right)$<br>It results<br>$\left\|\tilde{G}_{28}^{(1)}-\tilde{G}_i^{(2)}\right\|\leq\int_0^t(a_{28})^{(5)}\left\|G_{29}^{(1)}-G_{29}^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{(\widehat{M}_{28})^{(5)}s_{(28)}}ds_{(28)}+$<br>$\int_0^t\{(a_{28}')^{(5)}\left\|G_{28}^{(1)}-G_{28}^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}+$<br>$(a_{28}'')^{(5)}\left(T_{29}^{(1)},s_{(28)}\right)\left\|G_{28}^{(1)}-G_{28}^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{(\widehat{M}_{28})^{(5)}s_{(28)}}+$<br>$G_{28}^{(2)}\left\|(a_{28}'')^{(5)}\left(T_{29}^{(1)},s_{(28)}\right)-(a_{28}'')^{(5)}\left(T_{29}^{(2)},s_{(28)}\right)\right\|\ e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{(\widehat{M}_{28})^{(5)}s_{(28)}}\}ds_{(28)}$<br>Where $s_{(28)}$ represents integrand that is integrated over the interval $[0,t]$<br>From the hypotheses on it follows | 236 |
| $\left\|(G_{31})^{(1)}-(G_{31})^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}t}\leq$<br>$\frac{1}{(\widehat{M}_{28})^{(5)}}\left((a_{28})^{(5)}+(a_{28}')^{(5)}+(\widehat{A}_{28})^{(5)}+(\widehat{P}_{28})^{(5)}(\widehat{k}_{28})^{(5)}\right)d\left(((G_{31})^{(1)},(T_{31})^{(1)};\ (G_{31})^{(2)},(T_{31})^{(2)})\right)$<br>And analogous inequalities for $G_i\ and\ T_i$. Taking into account the hypothesis the result follows | 237 |
| **Remark 21:** The fact that we supposed $(a_{28}'')^{(5)}$ and $(b_{28}'')^{(5)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{28})^{(5)}e^{(\widehat{M}_{28})^{(5)}t}\ and\ (\widehat{Q}_{28})^{(5)}e^{(\widehat{M}_{28})^{(5)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it | 238 |

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| suffices to consider that $(a_i^{''})^{(5)}$ and $(b_i^{''})^{(5)}, i = 28,29,30$ depend only on $T_{29}$ and respectively on $(G_{31})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 22:** There does not exist any $t$ where $G_i\,(t) = 0\ and\ T_i\,(t) = 0$<br><br>it results<br><br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(5)} - (a_i^{''})^{(5)}(T_{29}(s_{(28)}), s_{(28)})\} ds_{(28)}\right]} \geq 0$<br><br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(5)} t\right)} > 0$ for t > 0 | 239 |
| **Definition of** $\left((\widehat{M}_{28})^{(5)}\right)_1, \left((\widehat{M}_{28})^{(5)}\right)_2\ and\ \left((\widehat{M}_{28})^{(5)}\right)_3$ :<br><br>**Remark 23:** if $G_{28}$ is bounded, the same property have also $G_{29}\ and\ G_{30}$ . indeed if<br><br>$G_{28} < (\widehat{M}_{28})^{(5)}$ it follows $\frac{dG_{29}}{dt} \leq \left((\widehat{M}_{28})^{(5)}\right)_1 - (a_{29}^{'})^{(5)} G_{29}$ and by integrating<br><br>$G_{29} \leq \left((\widehat{M}_{28})^{(5)}\right)_2 = G_{29}^0 + 2(a_{29})^{(5)}\left((\widehat{M}_{28})^{(5)}\right)_1 / (a_{29}^{'})^{(5)}$<br><br>In the same way , one can obtain<br><br>$G_{30} \leq \left((\widehat{M}_{28})^{(5)}\right)_3 = G_{30}^0 + 2(a_{30})^{(5)}\left((\widehat{M}_{28})^{(5)}\right)_2 / (a_{30}^{'})^{(5)}$<br><br>If $G_{29}\ or\ G_{30}$ is bounded, the same property follows for $G_{28}$ , $G_{30}$ and $G_{28}$ , $G_{29}$ respectively. | 240 |
| **Remark 24:** If $G_{28}$ *is* bounded, from below, the same property holds for $G_{29}\ and\ G_{30}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{29}$ is bounded from below. | 241 |
| **Remark 25:** If $T_{28}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(5)}((G_{31})(t), t)) = (b_{29}^{'})^{(5)}$ then $T_{29} \to \infty$.<br><br>**Definition of** $(m)^{(5)}$ and $\varepsilon_5$ :<br><br>Indeed let $t_5$ be so that for $t > t_5$<br><br>$(b_{29})^{(5)} - (b_i^{''})^{(5)}((G_{31})(t), t) < \varepsilon_5, T_{28}\,(t) > (m)^{(5)}$ | 242 |
| Then $\frac{dT_{29}}{dt} \geq (a_{29})^{(5)}(m)^{(5)} - \varepsilon_5 T_{29}$ which leads to<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{\varepsilon_5}\right)(1 - e^{-\varepsilon_5 t}) + T_{29}^0 e^{-\varepsilon_5 t}$ If we take t such that $e^{-\varepsilon_5 t} = \frac{1}{2}$ it results<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{2}\right), \quad t = log\frac{2}{\varepsilon_5}$ By taking now $\varepsilon_5$ sufficiently small one sees that $T_{29}$ is unbounded.<br>The same property holds for $T_{30}$ if $\lim_{t\to\infty}(b_{30}^{''})^{(5)}\left((G_{31})(t), t\right) = (b_{30}^{'})^{(5)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations<br><br>Analogous inequalities hold also for $G_{33}$ , $G_{34}, T_{32}, T_{33}, T_{34}$ | 243 |

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| It is now sufficient to take $\frac{(a_i)^{(6)}}{(\widehat{M}_{32})^{(6)}}$ , $\frac{(b_i)^{(6)}}{(\widehat{M}_{32})^{(6)}} < 1$ and to choose<br>$(\widehat{P}_{32})^{(6)}$ and $(\widehat{Q}_{32})^{(6)}$ large to have | 244 |
| $\frac{(a_i)^{(6)}}{(\widehat{M}_{32})^{(6)}}\left[(\widehat{P}_{32})^{(6)} + \left((\widehat{P}_{32})^{(6)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{32})^{(6)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{32})^{(6)}$ | 245 |
| $\frac{(b_i)^{(6)}}{(\widehat{M}_{32})^{(6)}}\left[\left((\widehat{Q}_{32})^{(6)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{32})^{(6)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{32})^{(6)}\right] \leq (\widehat{Q}_{32})^{(6)}$ | 246 |
| In order that the operator $\mathcal{A}^{(6)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(6)}$ is a contraction with respect to the metric<br>$d\left(((G_{35})^{(1)}, (T_{35})^{(1)}), ((G_{35})^{(2)}, (T_{35})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{32})^{(6)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{32})^{(6)}t}\}$<br>Indeed if we denote<br>**Definition of** $(\widetilde{G_{35}}), (\widetilde{T_{35}})$ : $\left((\widetilde{G_{35}}), (\widetilde{T_{35}})\right) = \mathcal{A}^{(6)}((G_{35}), (T_{35}))$<br>It results<br>$\left\|\tilde{G}_{32}^{(1)} - \tilde{G}_i^{(2)}\right\| \leq \int_0^t (a_{32})^{(6)}\left\|G_{33}^{(1)} - G_{33}^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{(\widehat{M}_{32})^{(6)}s_{(32)}}\,ds_{(32)} +$<br>$\int_0^t\{(a'_{32})^{(6)}\left\|G_{32}^{(1)} - G_{32}^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} +$<br>$(a''_{32})^{(6)}\left(T_{33}^{(1)}, s_{(32)}\right)\left\|G_{32}^{(1)} - G_{32}^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{(\widehat{M}_{32})^{(6)}s_{(32)}} +$<br>$G_{32}^{(2)}\left\|(a''_{32})^{(6)}\left(T_{33}^{(1)}, s_{(32)}\right) - (a''_{32})^{(6)}\left(T_{33}^{(2)}, s_{(32)}\right)\right\|\ e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{(\widehat{M}_{32})^{(6)}s_{(32)}}\}ds_{(32)}$<br>Where $s_{(32)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | 247 |
| $\left\|(G_{35})^{(1)} - (G_{35})^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}t} \leq$<br>$\frac{1}{(\widehat{M}_{32})^{(6)}}\left((a_{32})^{(6)} + (a'_{32})^{(6)} + (\widehat{A}_{32})^{(6)} + (\widehat{P}_{32})^{(6)}(\widehat{k}_{32})^{(6)}\right)d\left(((G_{35})^{(1)}, (T_{35})^{(1)}; (G_{35})^{(2)}, (T_{35})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 248 |
| **Remark 26:** The fact that we supposed $(a''_{32})^{(6)}$ and $(b''_{32})^{(6)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate | 249 |

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| condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{32})^{(6)}e^{(\widehat{M}_{32})^{(6)}t}$ $and$ $(\widehat{Q}_{32})^{(6)}e^{(\widehat{M}_{32})^{(6)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(6)}$ and $(b_i^{''})^{(6)}, i = 32,33,34$ depend only on $T_{33}$ and respectively on $(G_{35})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 27:** There does not exist any $t$ where $G_i(t) = 0\ and\ T_i(t) = 0$<br><br>it results<br><br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(6)} - (a_i^{''})^{(6)}(T_{33}(s_{(32)}), s_{(32)})\} ds_{(32)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{(-(b_i^{'})^{(6)}t)} > 0$ for t > 0 | 250 |
| **Definition of** $\left((\widehat{M}_{32})^{(6)}\right)_1, \left((\widehat{M}_{32})^{(6)}\right)_2$ $and$ $\left((\widehat{M}_{32})^{(6)}\right)_3$ :<br><br>**Remark 28:** if $G_{32}$ is bounded, the same property have also $G_{33}$ $and$ $G_{34}$ . indeed if<br><br>$G_{32} < (\widehat{M}_{32})^{(6)}$ it follows $\frac{dG_{33}}{dt} \leq \left((\widehat{M}_{32})^{(6)}\right)_1 - (a_{33}^{'})^{(6)} G_{33}$ and by integrating<br><br>$G_{33} \leq \left((\widehat{M}_{32})^{(6)}\right)_2 = G_{33}^0 + 2(a_{33})^{(6)}\left((\widehat{M}_{32})^{(6)}\right)_1 / (a_{33}^{'})^{(6)}$<br><br>In the same way , one can obtain<br><br>$G_{34} \leq \left((\widehat{M}_{32})^{(6)}\right)_3 = G_{34}^0 + 2(a_{34})^{(6)}\left((\widehat{M}_{32})^{(6)}\right)_2 / (a_{34}^{'})^{(6)}$<br><br>If $G_{33}$ $or$ $G_{34}$ is bounded, the same property follows for $G_{32}$ , $G_{34}$ and $G_{32}$ , $G_{33}$ respectively. | 251 |
| **Remark 29:** If $G_{32}$ $is$ bounded, from below, the same property holds for $G_{33}$ $and$ $G_{34}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{33}$ is bounded from below. | 252 |
| **Remark 30:** If $T_{32}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(6)}((G_{35})(t), t)) = (b_{33}^{'})^{(6)}$ then $T_{33} \to \infty$.<br><br>**Definition of** $(m)^{(6)}$ and $\varepsilon_6$ :<br><br>Indeed let $t_6$ be so that for $t > t_6$<br><br>$(b_{33})^{(6)} - (b_i^{''})^{(6)}((G_{35})(t), t) < \varepsilon_6, T_{32}(t) > (m)^{(6)}$ | 253 |
| Then $\frac{dT_{33}}{dt} \geq (a_{33})^{(6)}(m)^{(6)} - \varepsilon_6 T_{33}$ which leads to<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{\varepsilon_6}\right)(1 - e^{-\varepsilon_6 t}) + T_{33}^0 e^{-\varepsilon_6 t}$ If we take t such that $e^{-\varepsilon_6 t} = \frac{1}{2}$ it results<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{2}\right)$, $t = log\frac{2}{\varepsilon_6}$ By taking now $\varepsilon_6$ sufficiently small one sees that $T_{33}$ is unbounded.<br>The same property holds for $T_{34}$ if $\lim_{t\to\infty}(b_{34}^{''})^{(6)}\left((G_{35})(t), t(t), t\right) = (b_{34}^{'})^{(6)}$ | 254 |

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| We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |
| Analogous inequalities hold also for $G_{37}$ , $G_{38}, T_{36}, T_{37}, T_{38}$<br><br>It is now sufficient to take $\frac{(a_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}$ , $\frac{(b_i)^{(7)}}{(\widehat{M}_{36})^{(7)}} < 1$ and to choose<br>$(\widehat{\mathrm{P}}_{36})^{(7)}$ and $(\widehat{\mathrm{Q}}_{36})^{(7)}$ large to have | 255 |
| $\frac{(a_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}\left[(\widehat{P}_{36})^{(7)} + \left((\hat{P}_{36})^{(7)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{36})^{(7)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{36})^{(7)}$ | 256 |
| $\frac{(b_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}\left[\left((\hat{Q}_{36})^{(7)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{36})^{(7)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{36})^{(7)}\right] \leq (\hat{Q}_{36})^{(7)}$ | 257 |
| In order that the operator $\mathcal{A}^{(7)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(7)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{39})^{(1)}, (T_{39})^{(1)}), ((G_{39})^{(2)}, (T_{39})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{36})^{(7)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{36})^{(7)}t}\}$<br><br>Indeed if we denote<br><br>**Definition of** $(\widetilde{G_{39}}), (\widetilde{T_{39}}) : \left((\widetilde{G_{39}}), (\widetilde{T_{39}})\right) = \mathcal{A}^{(7)}((G_{39}), (T_{39}))$<br>It results<br><br>$\left\|\tilde{G}_{36}^{(1)} - \tilde{G}_i^{(2)}\right\| \leq \int_0^t (a_{36})^{(7)}\left\|G_{37}^{(1)} - G_{37}^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{(\widehat{M}_{36})^{(7)}s_{(36)}}\,ds_{(36)} +$<br>$\int_0^t\{(a'_{36})^{(7)}\left\|G_{36}^{(1)} - G_{36}^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{-(\widehat{M}_{36})^{(7)}s_{(36)}} +$<br>$(a''_{36})^{(7)}\left(T_{37}^{(1)}, s_{(36)}\right)\left\|G_{36}^{(1)} - G_{36}^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{(\widehat{M}_{36})^{(7)}s_{(36)}} +$<br>$G_{36}^{(2)}\left\|(a''_{36})^{(7)}\left(T_{37}^{(1)}, s_{(36)}\right) - (a''_{36})^{(7)}\left(T_{37}^{(2)}, s_{(36)}\right)\right\| \; e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{(\widehat{M}_{36})^{(7)}s_{(36)}}\}ds_{(36)}$<br><br>Where $s_{(36)}$ represents integrand that is integrated over the interval $[0, t]$<br><br>From the hypotheses on it follows | 258 |
| $\left\|(G_{39})^{(1)} - (G_{39})^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}t} \leq$<br>$\frac{1}{(\widehat{M}_{36})^{(7)}}\left((a_{36})^{(7)} + (a'_{36})^{(7)} + (\widehat{A}_{36})^{(7)} + (\widehat{P}_{36})^{(7)}(\widehat{k}_{36})^{(7)}\right)d\left(((G_{39})^{(1)}, (T_{39})^{(1)}; (G_{39})^{(2)}, (T_{39})^{(2)})\right)$<br><br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 259 |

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| **Remark 31:** The fact that we supposed $(a_{36}^{''})^{(7)}$ and $(b_{36}^{''})^{(7)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{36})^{(7)}e^{(\widehat{M}_{36})^{(7)}t}$ $and$ $(\widehat{Q}_{36})^{(7)}e^{(\widehat{M}_{36})^{(7)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(7)}$ and $(b_i^{''})^{(7)}, i = 36,37,38$ depend only on $T_{37}$ and respectively on $(G_{39})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 260 |
| **Remark 32:** There does not exist any $t$ where $G_i\,(t) = 0\ and\ T_i\,(t) = 0$<br><br>it results<br><br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t\{(a_i^{'})^{(7)}-(a_i^{''})^{(7)}(T_{37}(s_{(36)}),s_{(36)})\}ds_{(36)}\right]} \geq 0$<br><br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(7)}t\right)} > 0$ for t > 0 | 261 |
| **Definition of** $\left((\widehat{M}_{36})^{(7)}\right)_1, \left((\widehat{M}_{36})^{(7)}\right)_2\ and\ \left((\widehat{M}_{36})^{(7)}\right)_3$ :<br><br>**Remark 33:** if $G_{36}$ is bounded, the same property have also $G_{37}\ and\ G_{38}$ . indeed if<br><br>$G_{36} < (\widehat{M}_{36})^{(7)}$ it follows $\frac{dG_{37}}{dt} \leq \left((\widehat{M}_{36})^{(7)}\right)_1 - (a_{37}^{'})^{(7)}G_{37}$ and by integrating<br><br>$G_{37} \leq \left((\widehat{M}_{36})^{(7)}\right)_2 = G_{37}^0 + 2(a_{37})^{(7)}\left((\widehat{M}_{36})^{(7)}\right)_1/(a_{37}^{'})^{(7)}$<br><br>In the same way , one can obtain<br><br>$G_{38} \leq \left((\widehat{M}_{36})^{(7)}\right)_3 = G_{38}^0 + 2(a_{38})^{(7)}\left((\widehat{M}_{36})^{(7)}\right)_2/(a_{38}^{'})^{(7)}$<br><br>If $G_{37}\ or\ G_{38}$ is bounded, the same property follows for $G_{36}$ , $G_{38}$ and $G_{36}$ , $G_{37}$ respectively. | 262 |
| **Remark 34:** If $G_{36}$ $is$ bounded, from below, the same property holds for $G_{37}\ and\ G_{38}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{37}$ is bounded from below. | 263 |
| **Remark 35:** If $T_{36}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(7)}((G_{39})(t),t)) = (b_{37}^{'})^{(7)}$ then $T_{37} \to \infty$.<br><br>**Definition of** $(m)^{(7)}$ and $\varepsilon_7$ :<br><br>Indeed let $t_7$ be so that for $t > t_7$<br><br>$(b_{37})^{(7)} - (b_i^{''})^{(7)}((G_{39})(t),t) < \varepsilon_7, T_{36}\,(t) > (m)^{(7)}$ | 264 |

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| Then $\frac{dT_{37}}{dt} \geq (a_{37})^{(7)}(m)^{(7)} - \varepsilon_7 T_{37}$ which leads to<br><br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{\varepsilon_7}\right)(1 - e^{-\varepsilon_7 t}) + T_{37}^0 e^{-\varepsilon_7 t}$ If we take t such that $e^{-\varepsilon_7 t} = \frac{1}{2}$ it results<br><br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{2}\right), \quad t = log\frac{2}{\varepsilon_7}$ By taking now $\varepsilon_7$ sufficiently small one sees that $T_{37}$ is unbounded.<br>The same property holds for $T_{38}$ if $\lim_{t\to\infty}(b_{38}^{''})^{(7)}\left((G_{39})(t), t\right) = (b_{38}^{'})^{(7)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 265 |
| It is now sufficient to take $\frac{(a_i)^{(8)}}{(\bar{M}_{40})^{(8)}}$ , $\frac{(b_i)^{(8)}}{(\bar{M}_{40})^{(8)}} < 1$ and to choose<br>$(\hat{P}_{40})^{(8)}$ and $(\hat{Q}_{40})^{(8)}$ large to have | 266 |
| $\frac{(a_i)^{(8)}}{(\bar{M}_{40})^{(8)}}\left[(\hat{P}_{40})^{(8)} + \left((\hat{P}_{40})^{(8)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{40})^{(8)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{40})^{(8)}$ | 267 |
| $\frac{(b_i)^{(8)}}{(\bar{M}_{40})^{(8)}}\left[\left((\hat{Q}_{40})^{(8)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{40})^{(8)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{40})^{(8)}\right] \leq (\hat{Q}_{40})^{(8)}$ | 268 |
| In order that the operator $\mathcal{A}^{(8)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(8)}$ is a contraction with respect to the metric | |
| $d\left(((G_{43})^{(1)}, (T_{43})^{(1)}), ((G_{43})^{(2)}, (T_{43})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\bar{M}_{40})^{(8)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\bar{M}_{40})^{(8)}t}\}$ | 269 |
| Indeed if we denote<br><br>**Definition of** $\widetilde{(G_{43})}, \widetilde{(T_{43})}$ : $\left(\widetilde{(G_{43})}, \widetilde{(T_{43})}\right) = \mathcal{A}^{(8)}((G_{43}), (T_{43}))$ | 270 |
| It results<br><br>$\left\|\tilde{G}_{40}^{(1)} - \tilde{G}_i^{(2)}\right\| \leq \int_0^t (a_{40})^{(8)}\left\|G_{41}^{(1)} - G_{41}^{(2)}\right\|e^{-(\bar{M}_{40})^{(8)}s_{(40)}}e^{(\bar{M}_{40})^{(8)}s_{(40)}}\,ds_{(40)} +$<br>$\int_0^t\{(a_{40}^{'})^{(8)}\left\|G_{40}^{(1)} - G_{40}^{(2)}\right\|e^{-(\bar{M}_{40})^{(8)}s_{(40)}}e^{-(\bar{M}_{40})^{(8)}s_{(40)}} +$<br>$(a_{40}^{''})^{(8)}\left(T_{41}^{(1)}, s_{(40)}\right)\left\|G_{40}^{(1)} - G_{40}^{(2)}\right\|e^{-(\bar{M}_{40})^{(8)}s_{(40)}}e^{(\bar{M}_{40})^{(8)}s_{(40)}} +$<br>$G_{40}^{(2)}\left\|(a_{40}^{''})^{(8)}\left(T_{41}^{(1)}, s_{(40)}\right) - (a_{40}^{''})^{(8)}\left(T_{41}^{(2)}, s_{(40)}\right)\right\| e^{-(\bar{M}_{40})^{(8)}s_{(40)}}e^{(\bar{M}_{40})^{(8)}s_{(40)}}\}ds_{(40)}$ | 271 |
| Where $s_{(40)}$ represents integrand that is integrated over the interval [0, t] | 272 |

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| From the hypotheses it follows | |
| $\left\|(G_{43})^{(1)} - (G_{43})^{(2)}\right\| e^{-(\widehat{M}_{40})^{(8)}t} \leq$ $\frac{1}{(\widehat{M}_{40})^{(8)}}\left((a_{40})^{(8)} + (a'_{40})^{(8)} + (\widehat{A}_{40})^{(8)} + (\widehat{P}_{40})^{(8)}(\widehat{k}_{40})^{(8)}\right) d\left(((G_{43})^{(1)}, (T_{43})^{(1)};\ (G_{43})^{(2)}, (T_{43})^{(2)})\right)$ And analogous inequalities for $G_i$ $and$ $T_i$. Taking into account the hypothesis the result follows | 273 |
| **Remark 36:** The fact that we supposed $(a''_{40})^{(8)}$ and $(b''_{40})^{(8)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{40})^{(8)} e^{(\widehat{M}_{40})^{(8)}t}$ $and$ $(\widehat{Q}_{40})^{(8)} e^{(\widehat{M}_{40})^{(8)}t}$ respectively of $\mathbb{R}_+$. If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a''_i)^{(8)}$ and $(b''_i)^{(8)}, i = 40,41,42$ depend only on $T_{41}$ and respectively on $(G_{43})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 274 |
| **Remark 37** There does not exist any $t$ where $G_i(t) = 0$ $and$ $T_i(t) = 0$ it results $G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a'_i)^{(8)} - (a''_i)^{(8)}(T_{41}(s_{(40)}), s_{(40)})\} ds_{(40)}\right]} \geq 0$ $T_i(t) \geq T_i^0 e^{\left(-(b'_i)^{(8)}t\right)} > 0$ for t > 0 | 275 |
| **Definition of** $\left((\widehat{M}_{40})^{(8)}\right)_1, \left((\widehat{M}_{40})^{(8)}\right)_2$ $and$ $\left((\widehat{M}_{40})^{(8)}\right)_3$ : **Remark 38:** if $G_{40}$ is bounded, the same property have also $G_{41}$ $and$ $G_{42}$ . indeed if $G_{40} < (\widehat{M}_{40})^{(8)}$ it follows $\frac{dG_{41}}{dt} \leq \left((\widehat{M}_{40})^{(8)}\right)_1 - (a'_{41})^{(8)} G_{41}$ and by integrating $G_{41} \leq \left((\widehat{M}_{40})^{(8)}\right)_2 = G_{41}^0 + 2(a_{41})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_1 / (a'_{41})^{(8)}$ In the same way , one can obtain $G_{42} \leq \left((\widehat{M}_{40})^{(8)}\right)_3 = G_{42}^0 + 2(a_{42})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_2 / (a'_{42})^{(8)}$ If $G_{41}$ $or$ $G_{42}$ is bounded, the same property follows for $G_{40}$ , $G_{42}$ and $G_{40}$ , $G_{41}$ respectively. | 276 |
| **Remark 39:** If $G_{40}$ $is$ bounded, from below, the same property holds for $G_{41}$ $and$ $G_{42}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{41}$ is bounded from below. | 277 |
| **Remark 40:** If $T_{40}$ is bounded from below and $\lim_{t\to\infty}((b''_i)^{(8)}((G_{43})(t), t)) = (b'_{41})^{(8)}$ then $T_{41} \to \infty$. | 278 |

**Definition of** $(m)^{(8)}$ and $\varepsilon_8$ :

Indeed let $t_8$ be so that for $t > t_8$

$(b_{41})^{(8)} - (b_i^{''})^{(8)}\big((G_{43})(t), t\big) < \varepsilon_8, T_{40}(t) > (m)^{(8)}$

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Then $\frac{dT_{41}}{dt} \geq (a_{41})^{(8)}(m)^{(8)} - \varepsilon_8 T_{41}$ which leads to

$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{\varepsilon_8}\right)(1 - e^{-\varepsilon_8 t}) + T_{41}^0 e^{-\varepsilon_8 t}$ If we take t such that $e^{-\varepsilon_8 t} = \frac{1}{2}$ it results

$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{2}\right), \quad t = log\frac{2}{\varepsilon_8}$ By taking now $\varepsilon_8$ sufficiently small one sees that $T_{41}$ is unbounded.
The same property holds for $T_{42}$ if $\lim_{t\to\infty}(b_{42}^{''})^{(8)}\big((G_{43})(t), t(t), t\big) = (b_{42}^{'})^{(8)}$

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It is now sufficient to take $\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}, \frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}} < 1$ and to choose $(\widehat{P}_{44})^{(9)}$ and $(\widehat{Q}_{44})^{(9)}$ large to have

$$\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[(\widehat{P}_{44})^{(9)} + \big((\hat{P}_{44})^{(9)} + G_j^0\big)e^{-\left(\frac{(\hat{P}_{44})^{(9)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{44})^{(9)}$$

$$\frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[\big((\hat{Q}_{44})^{(9)} + T_j^0\big)e^{-\left(\frac{(\hat{Q}_{44})^{(9)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{44})^{(9)}\right] \leq (\hat{Q}_{44})^{(9)}$$

In order that the operator $\mathcal{A}^{(9)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying 39,35,36 into itself

The operator $\mathcal{A}^{(9)}$ is a contraction with respect to the metric

$$d\Big(\big((G_{47})^{(1)}, (T_{47})^{(1)}\big), \big((G_{47})^{(2)}, (T_{47})^{(2)}\big)\Big) =$$
$$\sup_i\{\max_{t\in\mathbb{R}_+}\big|G_i^{(1)}(t) - G_i^{(2)}(t)\big|e^{-(\widehat{M}_{44})^{(9)}t}, \max_{t\in\mathbb{R}_+}\big|T_i^{(1)}(t) - T_i^{(2)}(t)\big|e^{-(\widehat{M}_{44})^{(9)}t}\}$$

Indeed if we denote

**Definition of** $\widetilde{(G_{47})}, \widetilde{(T_{47})} : \big(\widetilde{(G_{47})}, \widetilde{(T_{47})}\big) = \mathcal{A}^{(9)}\big((G_{47}), (T_{47})\big)$
It results

$$\big|\tilde{G}_{44}^{(1)} - \tilde{G}_i^{(2)}\big| \leq \int_0^t (a_{44})^{(9)}\big|G_{45}^{(1)} - G_{45}^{(2)}\big|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\,ds_{(44)} +$$
$$\int_0^t\{(a_{44}^{'})^{(9)}\big|G_{44}^{(1)} - G_{44}^{(2)}\big|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{-(\widehat{M}_{44})^{(9)}s_{(44)}} +$$
$$(a_{44}^{''})^{(9)}\big(T_{45}^{(1)}, s_{(44)}\big)\big|G_{44}^{(1)} - G_{44}^{(2)}\big|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}} +$$
$$G_{44}^{(2)}\big|(a_{44}^{''})^{(9)}\big(T_{45}^{(1)}, s_{(44)}\big) - (a_{44}^{''})^{(9)}\big(T_{45}^{(2)}, s_{(44)}\big)\big|\ e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\}ds_{(44)}$$

Where $s_{(44)}$ represents integrand that is integrated over the interval $[0, t]$

From the hypotheses on 45,46,47,28 and 29 it follows

$$\left|(G_{47})^{(1)} - G^{(2)}\right| e^{-(\widehat{M}_{44})^{(9)}t} \leq$$
$$\frac{1}{(\widehat{M}_{44})^{(9)}}\left((a_{44})^{(9)} + (a'_{44})^{(9)} + (\widehat{A}_{44})^{(9)} + (\widehat{P}_{44})^{(9)}(\widehat{k}_{44})^{(9)}\right) d\left(((G_{47})^{(1)}, (T_{47})^{(1)};\ (G_{47})^{(2)}, (T_{47})^{(2)})\right)$$

And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis (39,35,36) the result follows

**Remark 41:** The fact that we supposed $(a''_{44})^{(9)}$ and $(b''_{44})^{(9)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{44})^{(9)}e^{(\widehat{M}_{44})^{(9)}t}$ *and* $(\widehat{Q}_{44})^{(9)}e^{(\widehat{M}_{44})^{(9)}t}$ respectively of $\mathbb{R}_+$.

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i'')^{(9)}$ and $(b_i'')^{(9)}, i = 44,45,46$ depend only on $T_{45}$ and respectively on $(G_{47})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition.

**Remark 42:** There does not exist any $t$ where $G_i(t) = 0\ and\ T_i(t) = 0$

From 99 to 44 it results

$$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i')^{(9)} - (a_i'')^{(9)}(T_{45}(s_{(44)}), s_{(44)})\} ds_{(44)}\right]} \geq 0$$

$$T_i(t) \geq T_i^0 e^{(-(b_i')^{(9)}t)} > 0 \quad \text{for t} > 0$$

**Definition of** $\left((\widehat{M}_{44})^{(9)}\right)_1, \left((\widehat{M}_{44})^{(9)}\right)_2$ *and* $\left((\widehat{M}_{44})^{(9)}\right)_3$ :

**Remark 43:** if $G_{44}$ is bounded, the same property have also $G_{45}\ and\ G_{46}$ . indeed if

$G_{44} < (\widehat{M}_{44})^{(9)}$ it follows $\frac{dG_{45}}{dt} \leq \left((\widehat{M}_{44})^{(9)}\right)_1 - (a'_{45})^{(9)}G_{45}$ and by integrating

$$G_{45} \leq \left((\widehat{M}_{44})^{(9)}\right)_2 = G_{45}^0 + 2(a_{45})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_1/(a'_{45})^{(9)}$$

In the same way , one can obtain

$$G_{46} \leq \left((\widehat{M}_{44})^{(9)}\right)_3 = G_{46}^0 + 2(a_{46})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_2/(a'_{46})^{(9)}$$

If $G_{45}\ or\ G_{46}$ is bounded, the same property follows for $G_{44}$ , $G_{46}$ and $G_{44}$ , $G_{45}$ respectively.

**Remark 44:** If $G_{44}$ *is* bounded, from below, the same property holds for $G_{45}\ and\ G_{46}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{45}$ is bounded from below.

**Remark 45:** If $T_{44}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(9)}\left((G_{47})(t), t\right)) = (b'_{45})^{(9)}$ then $T_{45} \to \infty$.

**Definition of** $(m)^{(9)}$ and $\varepsilon_9$ :

Indeed let $t_9$ be so that for $t > t_9$

$$(b_{45})^{(9)} - (b_i^{''})^{(9)}\big((G_{47})(t),t\big) < \varepsilon_9, T_{44}\,(t) > (m)^{(9)}$$

Then $\frac{dT_{45}}{dt} \geq (a_{45})^{(9)}(m)^{(9)} - \varepsilon_9 T_{45}$ which leads to

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{\varepsilon_9}\right)(1 - e^{-\varepsilon_9 t}) + T_{45}^0 e^{-\varepsilon_9 t}$ If we take t such that $e^{-\varepsilon_9 t} = \frac{1}{2}$ it results

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{2}\right), \quad t = log\frac{2}{\varepsilon_9}$ By taking now $\varepsilon_9$ sufficiently small one sees that $T_{45}$ is unbounded.
The same property holds for $T_{46}$ if $\lim_{t\to\infty}(b_{46}^{''})^{(9)}\left((G_{47})(t),t\right) = (b_{46}^{'})^{(9)}$
We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 92

**Behavior of the solutions of equation** 280

**Theorem** If we denote and define

**Definition of** $(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ :

$(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ four constants satisfying
$-(\sigma_2)^{(1)} \leq -(a_{13}^{'})^{(1)} + (a_{14}^{'})^{(1)} - (a_{13}^{''})^{(1)}(T_{14}, t) + (a_{14}^{''})^{(1)}(T_{14}, t) \leq -(\sigma_1)^{(1)}$

$-(\tau_2)^{(1)} \leq -(b_{13}^{'})^{(1)} + (b_{14}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G,t) - (b_{14}^{''})^{(1)}(G,t) \leq -(\tau_1)^{(1)}$

**Definition of** $(\nu_1)^{(1)}, (\nu_2)^{(1)}, (u_1)^{(1)}, (u_2)^{(1)}, \nu^{(1)}, u^{(1)}$ : 281

By $(\nu_1)^{(1)} > 0$, $(\nu_2)^{(1)} < 0$ and respectively $(u_1)^{(1)} > 0$, $(u_2)^{(1)} < 0$ the roots of the equations
$(a_{14})^{(1)}\left(\nu^{(1)}\right)^2 + (\sigma_1)^{(1)}\nu^{(1)} - (a_{13})^{(1)} = 0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2 + (\tau_1)^{(1)}u^{(1)} - (b_{13})^{(1)} = 0$

**Definition of** $(\bar{\nu}_1)^{(1)},, (\bar{\nu}_2)^{(1)}, (\bar{u}_1)^{(1)}, (\bar{u}_2)^{(1)}$ : 282

By $(\bar{\nu}_1)^{(1)} > 0$, $(\bar{\nu}_2)^{(1)} < 0$ and respectively $(\bar{u}_1)^{(1)} > 0$, $(\bar{u}_2)^{(1)} < 0$ the roots of the equations
$(a_{14})^{(1)}\left(\nu^{(1)}\right)^2 + (\sigma_2)^{(1)}\nu^{(1)} - (a_{13})^{(1)} = 0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2 + (\tau_2)^{(1)}u^{(1)} - (b_{13})^{(1)} = 0$

**Definition of** $(m_1)^{(1)}, (m_2)^{(1)}, (\mu_1)^{(1)}, (\mu_2)^{(1)}, (\nu_0)^{(1)}$ :- 283

If we define $(m_1)^{(1)}, (m_2)^{(1)}, (\mu_1)^{(1)}, (\mu_2)^{(1)}$ by
$(m_2)^{(1)} = (\nu_0)^{(1)}, (m_1)^{(1)} = (\nu_1)^{(1)}, \; if\; (\nu_0)^{(1)} < (\nu_1)^{(1)}$

$(m_2)^{(1)} = (\nu_1)^{(1)}, (m_1)^{(1)} = (\bar{\nu}_1)^{(1)}, if\; (\nu_1)^{(1)} < (\nu_0)^{(1)} < (\bar{\nu}_1)^{(1)},$

and $\boxed{(\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}}$

$(m_2)^{(1)} = (\nu_1)^{(1)}, (m_1)^{(1)} = (\nu_0)^{(1)}, \; if\; (\bar{\nu}_1)^{(1)} < (\nu_0)^{(1)}$

and analogously 284

$(\mu_2)^{(1)} = (u_0)^{(1)}, (\mu_1)^{(1)} = (u_1)^{(1)}, \; if\; (u_0)^{(1)} < (u_1)^{(1)}$

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| $(\mu_2)^{(1)} = (u_1)^{(1)}, (\mu_1)^{(1)} = (\bar{u}_1)^{(1)}, if\ (u_1)^{(1)} < (u_0)^{(1)} < (\bar{u}_1)^{(1)},$ and $\boxed{(u_0)^{(1)} = \frac{T_{13}^0}{T_{14}^0}}$ $(\mu_2)^{(1)} = (u_1)^{(1)}, (\mu_1)^{(1)} = (u_0)^{(1)}, if\ (\bar{u}_1)^{(1)} < (u_0)^{(1)}$ where $(u_1)^{(1)}, (\bar{u}_1)^{(1)}$ are defined | |
| Then the solution of global equations satisfies the inequalities $G_{13}^0 e^{((S_1)^{(1)} - (p_{13})^{(1)})t} \le G_{13}(t) \le G_{13}^0 e^{(S_1)^{(1)} t}$ where $(p_i)^{(1)}$ is defined by equation $\frac{1}{(m_1)^{(1)}} G_{13}^0 e^{((S_1)^{(1)} - (p_{13})^{(1)})t} \le G_{14}(t) \le \frac{1}{(m_2)^{(1)}} G_{13}^0 e^{(S_1)^{(1)} t}$ | 285 |
| $\left(\frac{(a_{15})^{(1)} G_{13}^0}{(m_1)^{(1)}((S_1)^{(1)} - (p_{13})^{(1)} - (S_2)^{(1)})}\left[e^{((S_1)^{(1)} - (p_{13})^{(1)})t} - e^{-(S_2)^{(1)} t}\right] + G_{15}^0 e^{-(S_2)^{(1)} t} \le G_{15}(t) \le \frac{(a_{15})^{(1)} G_{13}^0}{(m_2)^{(1)}((S_1)^{(1)} - (a_{15}')^{(1)})}[e^{(S_1)^{(1)} t} - e^{-(a_{15}')^{(1)} t}] + G_{15}^0 e^{-(a_{15}')^{(1)} t}\right)$ | 286 |
| $\boxed{T_{13}^0 e^{(R_1)^{(1)} t} \le T_{13}(t) \le T_{13}^0 e^{((R_1)^{(1)} + (r_{13})^{(1)})t}}$ | 287 |
| $\frac{1}{(\mu_1)^{(1)}} T_{13}^0 e^{(R_1)^{(1)} t} \le T_{13}(t) \le \frac{1}{(\mu_2)^{(1)}} T_{13}^0 e^{((R_1)^{(1)} + (r_{13})^{(1)})t}$ | 288 |
| $\frac{(b_{15})^{(1)} T_{13}^0}{(\mu_1)^{(1)}((R_1)^{(1)} - (b_{15}')^{(1)})}\left[e^{(R_1)^{(1)} t} - e^{-(b_{15}')^{(1)} t}\right] + T_{15}^0 e^{-(b_{15}')^{(1)} t} \le T_{15}(t) \le \frac{(a_{15})^{(1)} T_{13}^0}{(\mu_2)^{(1)}((R_1)^{(1)} + (r_{13})^{(1)} + (R_2)^{(1)})}\left[e^{((R_1)^{(1)} + (r_{13})^{(1)})t} - e^{-(R_2)^{(1)} t}\right] + T_{15}^0 e^{-(R_2)^{(1)} t}$ | 289 |
| **Definition of** $(S_1)^{(1)}, (S_2)^{(1)}, (R_1)^{(1)}, (R_2)^{(1)}$:- Where $(S_1)^{(1)} = (a_{13})^{(1)}(m_2)^{(1)} - (a_{13}')^{(1)}$ $(S_2)^{(1)} = (a_{15})^{(1)} - (p_{15})^{(1)}$ $(R_1)^{(1)} = (b_{13})^{(1)}(\mu_2)^{(1)} - (b_{13}')^{(1)}$ $(R_2)^{(1)} = (b_{15}')^{(1)} - (r_{15})^{(1)}$ | 290 |
| **Behavior of the solutions of equation** **Theorem 2:** If we denote and define | 291 |
| **Definition of** $(\sigma_1)^{(2)}, (\sigma_2)^{(2)}, (\tau_1)^{(2)}, (\tau_2)^{(2)}$ : $(\sigma_1)^{(2)}, (\sigma_2)^{(2)}, (\tau_1)^{(2)}, (\tau_2)^{(2)}$ four constants satisfying | 292 |

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| $-(\sigma_2)^{(2)} \leq -(a_{16}^{'})^{(2)} + (a_{17}^{'})^{(2)} - (a_{16}^{''})^{(2)}(\mathrm{T}_{17}, t) + (a_{17}^{''})^{(2)}(\mathrm{T}_{17}, t) \leq -(\sigma_1)^{(2)}$ | 293 |
| $-(\tau_2)^{(2)} \leq -(b_{16}^{'})^{(2)} + (b_{17}^{'})^{(2)} - (b_{16}^{''})^{(2)}\big((G_{19}), t\big) - (b_{17}^{''})^{(2)}\big((G_{19}), t\big) \leq -(\tau_1)^{(2)}$ | 294 |
| **Definition of** $(\nu_1)^{(2)}, (\nu_2)^{(2)}, (u_1)^{(2)}, (u_2)^{(2)}$ : | 295 |
| By $(\nu_1)^{(2)} > 0$ , $(\nu_2)^{(2)} < 0$ and respectively $(u_1)^{(2)} > 0$ , $(u_2)^{(2)} < 0$ the roots | 296 |
| of the equations $(a_{17})^{(2)}\big(\nu^{(2)}\big)^2 + (\sigma_1)^{(2)}\nu^{(2)} - (a_{16})^{(2)} = 0$ | 297 |
| and $(b_{14})^{(2)}\big(u^{(2)}\big)^2 + (\tau_1)^{(2)}u^{(2)} - (b_{16})^{(2)} = 0$ and | 298 |
| **Definition of** $(\bar{\nu}_1)^{(2)}$,, $(\bar{\nu}_2)^{(2)}, (\bar{u}_1)^{(2)}, (\bar{u}_2)^{(2)}$ : | 299 |
| By $(\bar{\nu}_1)^{(2)} > 0$ , $(\bar{\nu}_2)^{(2)} < 0$ and respectively $(\bar{u}_1)^{(2)} > 0$ , $(\bar{u}_2)^{(2)} < 0$ the | 300 |
| roots of the equations $(a_{17})^{(2)}\big(\nu^{(2)}\big)^2 + (\sigma_2)^{(2)}\nu^{(2)} - (a_{16})^{(2)} = 0$ | 301 |
| and $(b_{17})^{(2)}\big(u^{(2)}\big)^2 + (\tau_2)^{(2)}u^{(2)} - (b_{16})^{(2)} = 0$ | 302 |
| **Definition of** $(m_1)^{(2)}$ , $(m_2)^{(2)}$ , $(\mu_1)^{(2)}, (\mu_2)^{(2)}$ :- | 303 |
| If we define $(m_1)^{(2)}$ , $(m_2)^{(2)}$ , $(\mu_1)^{(2)}, (\mu_2)^{(2)}$ by | 304 |
| $(m_2)^{(2)} = (\nu_0)^{(2)}, (m_1)^{(2)} = (\nu_1)^{(2)}$, ***if*** $(\nu_0)^{(2)} < (\nu_1)^{(2)}$ | 305 |
| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ , ***if*** $(\nu_1)^{(2)} < (\nu_0)^{(2)} < (\bar{\nu}_1)^{(2)}$, and $\boxed{(\nu_0)^{(2)} = \frac{\mathrm{G}_{16}^0}{\mathrm{G}_{17}^0}}$ | 306 |
| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\nu_0)^{(2)}$, ***if*** $(\bar{\nu}_1)^{(2)} < (\nu_0)^{(2)}$ | 307 |
| and analogously<br>$(\mu_2)^{(2)} = (u_0)^{(2)}, (\mu_1)^{(2)} = (u_1)^{(2)}$, ***if*** $(u_0)^{(2)} < (u_1)^{(2)}$<br>$(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (\bar{u}_1)^{(2)}$ , ***if*** $(u_1)^{(2)} < (u_0)^{(2)} < (\bar{u}_1)^{(2)}$,<br>and $\boxed{(u_0)^{(2)} = \frac{\mathrm{T}_{16}^0}{\mathrm{T}_{17}^0}}$ | 308 |
| $(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (u_0)^{(2)}$, ***if*** $(\bar{u}_1)^{(2)} < (u_0)^{(2)}$ | 309 |
| Then the solution of global equations satisfies the inequalities<br>$\mathrm{G}_{16}^0 \mathrm{e}^{((\mathrm{S}_1)^{(2)} - (p_{16})^{(2)})\mathrm{t}} \leq G_{16}(t) \leq \mathrm{G}_{16}^0 \mathrm{e}^{(\mathrm{S}_1)^{(2)}\mathrm{t}}$ | 310 |
| $(p_i)^{(2)}$ is defined by equation | |

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| $\frac{1}{(m_1)^{(2)}} \mathrm{G}_{16}^{0} \mathrm{e}^{((\mathrm{S}_1)^{(2)}-(p_{16})^{(2)})\mathrm{t}} \leq G_{17}(t) \leq \frac{1}{(m_2)^{(2)}} \mathrm{G}_{16}^{0} \mathrm{e}^{(\mathrm{S}_1)^{(2)}\mathrm{t}}$ | 311 |
| $\left(\frac{(a_{18})^{(2)}\mathrm{G}_{16}^{0}}{(m_1)^{(2)}((\mathrm{S}_1)^{(2)}-(p_{16})^{(2)}-(\mathrm{S}_2)^{(2)})}\left[\mathrm{e}^{((\mathrm{S}_1)^{(2)}-(p_{16})^{(2)})\mathrm{t}} - \mathrm{e}^{-(\mathrm{S}_2)^{(2)}\mathrm{t}}\right] + \mathrm{G}_{18}^{0}\mathrm{e}^{-(\mathrm{S}_2)^{(2)}\mathrm{t}} \leq \mathrm{G}_{18}(t) \leq \right.$ $\left.\frac{(a_{18})^{(2)}\mathrm{G}_{16}^{0}}{(m_2)^{(2)}((\mathrm{S}_1)^{(2)}-(a_{18}^{'})^{(2)})}\left[\mathrm{e}^{(\mathrm{S}_1)^{(2)}\mathrm{t}} - \mathrm{e}^{-(a_{18}^{'})^{(2)}\mathrm{t}}\right] + \mathrm{G}_{18}^{0}\mathrm{e}^{-(a_{18}^{'})^{(2)}\mathrm{t}}\right)$ | 312 |
| $\boxed{\mathrm{T}_{16}^{0}\mathrm{e}^{(\mathrm{R}_1)^{(2)}t} \leq T_{16}(t) \leq \mathrm{T}_{16}^{0}\mathrm{e}^{((\mathrm{R}_1)^{(2)}+(r_{16})^{(2)})t}}$ | 313 |
| $\frac{1}{(\mu_1)^{(2)}} \mathrm{T}_{16}^{0}\mathrm{e}^{(\mathrm{R}_1)^{(2)}t} \leq T_{16}(t) \leq \frac{1}{(\mu_2)^{(2)}} \mathrm{T}_{16}^{0}\mathrm{e}^{((\mathrm{R}_1)^{(2)}+(r_{16})^{(2)})t}$ | 314 |
| $\frac{(b_{18})^{(2)}\mathrm{T}_{16}^{0}}{(\mu_1)^{(2)}((\mathrm{R}_1)^{(2)}-(b_{18}^{'})^{(2)})}\left[\mathrm{e}^{(\mathrm{R}_1)^{(2)}t} - \mathrm{e}^{-(b_{18}^{'})^{(2)}t}\right] + \mathrm{T}_{18}^{0}\mathrm{e}^{-(b_{18}^{'})^{(2)}t} \leq T_{18}(t) \leq$ $\frac{(a_{18})^{(2)}\mathrm{T}_{16}^{0}}{(\mu_2)^{(2)}((\mathrm{R}_1)^{(2)}+(r_{16})^{(2)}+(\mathrm{R}_2)^{(2)})}\left[\mathrm{e}^{((\mathrm{R}_1)^{(2)}+(r_{16})^{(2)})t} - \mathrm{e}^{-(\mathrm{R}_2)^{(2)}t}\right] + \mathrm{T}_{18}^{0}\mathrm{e}^{-(\mathrm{R}_2)^{(2)}t}$ | 315 |
| **Definition of** $(\mathrm{S}_1)^{(2)}, (\mathrm{S}_2)^{(2)}, (\mathrm{R}_1)^{(2)}, (\mathrm{R}_2)^{(2)}$:- | 316 |
| Where $(\mathrm{S}_1)^{(2)} = (a_{16})^{(2)}(m_2)^{(2)} - (a_{16}^{'})^{(2)}$ $(\mathrm{S}_2)^{(2)} = (a_{18})^{(2)} - (p_{18})^{(2)}$ | 317 |
| $(R_1)^{(2)} = (b_{16})^{(2)}(\mu_2)^{(1)} - (b_{16}^{'})^{(2)}$ $(\mathrm{R}_2)^{(2)} = (b_{18}^{'})^{(2)} - (r_{18})^{(2)}$ | 318 |
| **Behavior of the solutions** **Theorem 3:** If we denote and define **Definition of** $(\sigma_1)^{(3)}, (\sigma_2)^{(3)}, (\tau_1)^{(3)}, (\tau_2)^{(3)}$: $(\sigma_1)^{(3)}, (\sigma_2)^{(3)}, (\tau_1)^{(3)}, (\tau_2)^{(3)}$ four constants satisfying $-(\sigma_2)^{(3)} \leq -(a_{20}^{'})^{(3)} + (a_{21}^{'})^{(3)} - (a_{20}^{''})^{(3)}(T_{21}, t) + (a_{21}^{''})^{(3)}(T_{21}, t) \leq -(\sigma_1)^{(3)}$ $-(\tau_2)^{(3)} \leq -(b_{20}^{'})^{(3)} + (b_{21}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23}, t) - (b_{21}^{''})^{(3)}((G_{23}), t) \leq -(\tau_1)^{(3)}$ | 319 |
| **Definition of** $(\nu_1)^{(3)}, (\nu_2)^{(3)}, (u_1)^{(3)}, (u_2)^{(3)}$: By $(\nu_1)^{(3)} > 0$, $(\nu_2)^{(3)} < 0$ and respectively $(u_1)^{(3)} > 0$, $(u_2)^{(3)} < 0$ the roots of the equations $(a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_1)^{(3)}\nu^{(3)} - (a_{20})^{(3)} = 0$ and $(b_{21})^{(3)}\left(u^{(3)}\right)^2 + (\tau_1)^{(3)}u^{(3)} - (b_{20})^{(3)} = 0$ and By $(\bar{\nu}_1)^{(3)} > 0$, $(\bar{\nu}_2)^{(3)} < 0$ and respectively $(\bar{u}_1)^{(3)} > 0$, $(\bar{u}_2)^{(3)} < 0$ the | 320 |

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| roots of the equations $(a_{21})^{(3)}\left(v^{(3)}\right)^2 + (\sigma_2)^{(3)}v^{(3)} - (a_{20})^{(3)} = 0$<br><br>and $(b_{21})^{(3)}\left(u^{(3)}\right)^2 + (\tau_2)^{(3)}u^{(3)} - (b_{20})^{(3)} = 0$ | |
| **Definition of** $(m_1)^{(3)}, (m_2)^{(3)}, (\mu_1)^{(3)}, (\mu_2)^{(3)}$ :-<br><br>If we define $(m_1)^{(3)}, (m_2)^{(3)}, (\mu_1)^{(3)}, (\mu_2)^{(3)}$ by<br>$(m_2)^{(3)} = (v_0)^{(3)}, (m_1)^{(3)} = (v_1)^{(3)}$, ***if*** $(v_0)^{(3)} < (v_1)^{(3)}$<br><br>$(m_2)^{(3)} = (v_1)^{(3)}, (m_1)^{(3)} = (\bar{v}_1)^{(3)}$, ***if*** $(v_1)^{(3)} < (v_0)^{(3)} < (\bar{v}_1)^{(3)}$,<br><br>and $\boxed{(v_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0}}$<br><br>$(m_2)^{(3)} = (v_1)^{(3)}, (m_1)^{(3)} = (v_0)^{(3)}$, ***if*** $(\bar{v}_1)^{(3)} < (v_0)^{(3)}$ | 321 |
| and analogously<br><br>$(\mu_2)^{(3)} = (u_0)^{(3)}, (\mu_1)^{(3)} = (u_1)^{(3)}$, ***if*** $(u_0)^{(3)} < (u_1)^{(3)}$<br><br>$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (\bar{u}_1)^{(3)}$, ***if*** $(u_1)^{(3)} < (u_0)^{(3)} < (\bar{u}_1)^{(3)}$, and $\boxed{(u_0)^{(3)} = \frac{T_{20}^0}{T_{21}^0}}$<br><br>$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (u_0)^{(3)}$, ***if*** $(\bar{u}_1)^{(3)} < (u_0)^{(3)}$<br><br>Then the solution of global equations satisfies the inequalities<br><br>$G_{20}^0 e^{((S_1)^{(3)} - (p_{20})^{(3)})t} \le G_{20}(t) \le G_{20}^0 e^{(S_1)^{(3)}t}$<br><br>$(p_i)^{(3)}$ is defined by equation | 322 |
| $\frac{1}{(m_1)^{(3)}} G_{20}^0 e^{((S_1)^{(3)} - (p_{20})^{(3)})t} \le G_{21}(t) \le \frac{1}{(m_2)^{(3)}} G_{20}^0 e^{(S_1)^{(3)}t}$ | 323 |
| $\left(\frac{(a_{22})^{(3)} G_{20}^0}{(m_1)^{(3)}((S_1)^{(3)} - (p_{20})^{(3)} - (S_2)^{(3)})}\left[e^{((S_1)^{(3)} - (p_{20})^{(3)})t} - e^{-(S_2)^{(3)}t}\right] + G_{22}^0 e^{-(S_2)^{(3)}t} \le G_{22}(t) \le \right.$<br>$\left.\frac{(a_{22})^{(3)} G_{20}^0}{(m_2)^{(3)}((S_1)^{(3)} - (a'_{22})^{(3)})}\left[e^{(S_1)^{(3)}t} - e^{-(a'_{22})^{(3)}t}\right] + G_{22}^0 e^{-(a'_{22})^{(3)}t}\right)$ | 324 |
| $\boxed{T_{20}^0 e^{(R_1)^{(3)}t} \le T_{20}(t) \le T_{20}^0 e^{((R_1)^{(3)} + (r_{20})^{(3)})t}}$ | 325 |
| $\frac{1}{(\mu_1)^{(3)}} T_{20}^0 e^{(R_1)^{(3)}t} \le T_{20}(t) \le \frac{1}{(\mu_2)^{(3)}} T_{20}^0 e^{((R_1)^{(3)} + (r_{20})^{(3)})t}$ | 326 |
| $\frac{(b_{22})^{(3)} T_{20}^0}{(\mu_1)^{(3)}((R_1)^{(3)} - (b'_{22})^{(3)})}\left[e^{(R_1)^{(3)}t} - e^{-(b'_{22})^{(3)}t}\right] + T_{22}^0 e^{-(b'_{22})^{(3)}t} \le T_{22}(t) \le$<br><br>$\frac{(a_{22})^{(3)} T_{20}^0}{(\mu_2)^{(3)}((R_1)^{(3)} + (r_{20})^{(3)} + (R_2)^{(3)})}\left[e^{((R_1)^{(3)} + (r_{20})^{(3)})t} - e^{-(R_2)^{(3)}t}\right] + T_{22}^0 e^{-(R_2)^{(3)}t}$ | 327 |

**Definition of** $(S_1)^{(3)}, (S_2)^{(3)}, (R_1)^{(3)}, (R_2)^{(3)}$:- 328

Where $(S_1)^{(3)} = (a_{20})^{(3)}(m_2)^{(3)} - (a'_{20})^{(3)}$

$(S_2)^{(3)} = (a_{22})^{(3)} - (p_{22})^{(3)}$

$(R_1)^{(3)} = (b_{20})^{(3)}(\mu_2)^{(3)} - (b'_{20})^{(3)}$

$(R_2)^{(3)} = (b'_{22})^{(3)} - (r_{22})^{(3)}$

**Behavior of the solutions of equation**

**Theorem:** If we denote and define

**Definition of** $(\sigma_1)^{(4)}, (\sigma_2)^{(4)}, (\tau_1)^{(4)}, (\tau_2)^{(4)}$ :

$(\sigma_1)^{(4)}, (\sigma_2)^{(4)}, (\tau_1)^{(4)}, (\tau_2)^{(4)}$ four constants satisfying

$-(\sigma_2)^{(4)} \le -(a'_{24})^{(4)} + (a'_{25})^{(4)} - (a''_{24})^{(4)}(T_{25}, t) + (a''_{25})^{(4)}(T_{25}, t) \le -(\sigma_1)^{(4)}$

$-(\tau_2)^{(4)} \le -(b'_{24})^{(4)} + (b'_{25})^{(4)} - (b''_{24})^{(4)}((G_{27}), t) - (b''_{25})^{(4)}((G_{27}), t) \le -(\tau_1)^{(4)}$

**Definition of** $(v_1)^{(4)}, (v_2)^{(4)}, (u_1)^{(4)}, (u_2)^{(4)}, v^{(4)}, u^{(4)}$ : 329

By $(v_1)^{(4)} > 0$, $(v_2)^{(4)} < 0$ and respectively $(u_1)^{(4)} > 0$, $(u_2)^{(4)} < 0$ the roots of the equations $(a_{25})^{(4)}\left(v^{(4)}\right)^2 + (\sigma_1)^{(4)}v^{(4)} - (a_{24})^{(4)} = 0$ and $(b_{25})^{(4)}\left(u^{(4)}\right)^2 + (\tau_1)^{(4)}u^{(4)} - (b_{24})^{(4)} = 0$ and

**Definition of** $(\bar{v}_1)^{(4)}, , (\bar{v}_2)^{(4)}, (\bar{u}_1)^{(4)}, (\bar{u}_2)^{(4)}$ : 330

By $(\bar{v}_1)^{(4)} > 0$, $(\bar{v}_2)^{(4)} < 0$ and respectively $(\bar{u}_1)^{(4)} > 0$, $(\bar{u}_2)^{(4)} < 0$ the roots of the equations $(a_{25})^{(4)}\left(v^{(4)}\right)^2 + (\sigma_2)^{(4)}v^{(4)} - (a_{24})^{(4)} = 0$ and $(b_{25})^{(4)}\left(u^{(4)}\right)^2 + (\tau_2)^{(4)}u^{(4)} - (b_{24})^{(4)} = 0$

**Definition of** $(m_1)^{(4)}, (m_2)^{(4)}, (\mu_1)^{(4)}, (\mu_2)^{(4)}, (v_0)^{(4)}$ :-

If we define $(m_1)^{(4)}, (m_2)^{(4)}, (\mu_1)^{(4)}, (\mu_2)^{(4)}$ by

$(m_2)^{(4)} = (v_0)^{(4)}, (m_1)^{(4)} = (v_1)^{(4)}$, ***if*** $(v_0)^{(4)} < (v_1)^{(4)}$

$(m_2)^{(4)} = (v_1)^{(4)}, (m_1)^{(4)} = (\bar{v}_1)^{(4)}$, ***if*** $(v_4)^{(4)} < (v_0)^{(4)} < (\bar{v}_1)^{(4)}$,

and $\boxed{(v_0)^{(4)} = \frac{G^0_{24}}{G^0_{25}}}$

$(m_2)^{(4)} = (v_4)^{(4)}, (m_1)^{(4)} = (v_0)^{(4)}$, ***if*** $(\bar{v}_4)^{(4)} < (v_0)^{(4)}$

and analogously 331

$(\mu_2)^{(4)} = (u_0)^{(4)}, (\mu_1)^{(4)} = (u_1)^{(4)}$, ***if*** $(u_0)^{(4)} < (u_1)^{(4)}$

$(\mu_2)^{(4)} = (u_1)^{(4)}, (\mu_1)^{(4)} = (\bar{u}_1)^{(4)}$, ***if*** $(u_1)^{(4)} < (u_0)^{(4)} < (\bar{u}_1)^{(4)}$,

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| and $\boxed{(u_0)^{(4)} = \frac{T_{24}^0}{T_{25}^0}}$<br>$(\mu_2)^{(4)} = (u_1)^{(4)}, (\mu_1)^{(4)} = (u_0)^{(4)}, \boldsymbol{if}\ (\bar{u}_1)^{(4)} < (u_0)^{(4)}$ where $(u_1)^{(4)}, (\bar{u}_1)^{(4)}$ | |
| Then the solution of global equations satisfies the inequalities<br>$G_{24}^0 e^{((S_1)^{(4)} - (p_{24})^{(4)})t} \le G_{24}(t) \le G_{24}^0 e^{(S_1)^{(4)} t}$<br>where $(p_i)^{(4)}$ is defined by equation | 332 |
| $\frac{1}{(m_1)^{(4)}} G_{24}^0 e^{((S_1)^{(4)} - (p_{24})^{(4)})t} \le G_{25}(t) \le \frac{1}{(m_2)^{(4)}} G_{24}^0 e^{(S_1)^{(4)} t}$ | 333 |
| $\left(\frac{(a_{26})^{(4)} G_{24}^0}{(m_1)^{(4)}((S_1)^{(4)} - (p_{24})^{(4)} - (S_2)^{(4)})}\right)\left[e^{((S_1)^{(4)} - (p_{24})^{(4)})t} - e^{-(S_2)^{(4)} t}\right] + G_{26}^0 e^{-(S_2)^{(4)} t} \le G_{26}(t) \le$<br>*(a26)4G240(m2)4(S1)4−(a26′)4e(S1)4t−e−(a26′)4t+ G260e−(a26′)4t* | 334 |
| $\boxed{T_{24}^0 e^{(R_1)^{(4)} t} \le T_{24}(t) \le T_{24}^0 e^{((R_1)^{(4)} + (r_{24})^{(4)})t}}$ | |
| $\frac{1}{(\mu_1)^{(4)}} T_{24}^0 e^{(R_1)^{(4)} t} \le T_{24}(t) \le \frac{1}{(\mu_2)^{(4)}} T_{24}^0 e^{((R_1)^{(4)} + (r_{24})^{(4)})t}$ | 335 |
| $\frac{(b_{26})^{(4)} T_{24}^0}{(\mu_1)^{(4)}((R_1)^{(4)} - (b_{26}')^{(4)})}\left[e^{(R_1)^{(4)} t} - e^{-(b_{26}')^{(4)} t}\right] + T_{26}^0 e^{-(b_{26}')^{(4)} t} \le T_{26}(t) \le$<br>$\frac{(a_{26})^{(4)} T_{24}^0}{(\mu_2)^{(4)}((R_1)^{(4)} + (r_{24})^{(4)} + (R_2)^{(4)})}\left[e^{((R_1)^{(4)} + (r_{24})^{(4)})t} - e^{-(R_2)^{(4)} t}\right] + T_{26}^0 e^{-(R_2)^{(4)} t}$ | 336 |
| **Definition of** $(S_1)^{(4)}, (S_2)^{(4)}, (R_1)^{(4)}, (R_2)^{(4)}$:-<br>Where $(S_1)^{(4)} = (a_{24})^{(4)}(m_2)^{(4)} - (a_{24}')^{(4)}$<br>$(S_2)^{(4)} = (a_{26})^{(4)} - (p_{26})^{(4)}$<br>$(R_1)^{(4)} = (b_{24})^{(4)}(\mu_2)^{(4)} - (b_{24}')^{(4)}$<br>$(R_2)^{(4)} = (b_{26}')^{(4)} - (r_{26})^{(4)}$ | 337 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ :<br>$(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ four constants satisfying<br>$-(\sigma_2)^{(5)} \le -(a_{28}')^{(5)} + (a_{29}')^{(5)} - (a_{28}'')^{(5)}(T_{29}, t) + (a_{29}'')^{(5)}(T_{29}, t) \le -(\sigma_1)^{(5)}$<br>$-(\tau_2)^{(5)} \le -(b_{28}')^{(5)} + (b_{29}')^{(5)} - (b_{28}'')^{(5)}((G_{31}), t) - (b_{29}'')^{(5)}((G_{31}), t) \le -(\tau_1)^{(5)}$ | 338 |

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| **Definition of** $(v_1)^{(5)}, (v_2)^{(5)}, (u_1)^{(5)}, (u_2)^{(5)}, v^{(5)}, u^{(5)}$ :<br><br>By $(v_1)^{(5)} > 0$ , $(v_2)^{(5)} < 0$ and respectively $(u_1)^{(5)} > 0$ , $(u_2)^{(5)} < 0$ the roots of the equations $(a_{29})^{(5)}\left(v^{(5)}\right)^2 + (\sigma_1)^{(5)}v^{(5)} - (a_{28})^{(5)} = 0$<br>and $(b_{29})^{(5)}\left(u^{(5)}\right)^2 + (\tau_1)^{(5)}u^{(5)} - (b_{28})^{(5)} = 0$ and | 339 |
| **Definition of** $(\bar{v}_1)^{(5)},, (\bar{v}_2)^{(5)}, (\bar{u}_1)^{(5)}, (\bar{u}_2)^{(5)}$ :<br><br>By $(\bar{v}_1)^{(5)} > 0$ , $(\bar{v}_2)^{(5)} < 0$ and respectively $(\bar{u}_1)^{(5)} > 0$ , $(\bar{u}_2)^{(5)} < 0$ the roots of the equations $(a_{29})^{(5)}\left(v^{(5)}\right)^2 + (\sigma_2)^{(5)}v^{(5)} - (a_{28})^{(5)} = 0$<br>and $(b_{29})^{(5)}\left(u^{(5)}\right)^2 + (\tau_2)^{(5)}u^{(5)} - (b_{28})^{(5)} = 0$<br>**Definition of** $(m_1)^{(5)}$ , $(m_2)^{(5)}$ , $(\mu_1)^{(5)}, (\mu_2)^{(5)}, (v_0)^{(5)}$ :-<br><br>If we define $(m_1)^{(5)}$ , $(m_2)^{(5)}$ , $(\mu_1)^{(5)}, (\mu_2)^{(5)}$ by<br><br>$(m_2)^{(5)} = (v_0)^{(5)}, (m_1)^{(5)} = (v_1)^{(5)}$, ***if*** $(v_0)^{(5)} < (v_1)^{(5)}$<br><br>$(m_2)^{(5)} = (v_1)^{(5)}, (m_1)^{(5)} = (\bar{v}_1)^{(5)}$ , ***if*** $(v_1)^{(5)} < (v_0)^{(5)} < (\bar{v}_1)^{(5)}$,<br>and $\boxed{(v_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}}$<br><br>$(m_2)^{(5)} = (v_1)^{(5)}, (m_1)^{(5)} = (v_0)^{(5)}$, ***if*** $(\bar{v}_1)^{(5)} < (v_0)^{(5)}$ | 340 |
| and analogously<br><br>$(\mu_2)^{(5)} = (u_0)^{(5)}, (\mu_1)^{(5)} = (u_1)^{(5)}$, ***if*** $(u_0)^{(5)} < (u_1)^{(5)}$<br><br>$(\mu_2)^{(5)} = (u_1)^{(5)}, (\mu_1)^{(5)} = (\bar{u}_1)^{(5)}$ , ***if*** $(u_1)^{(5)} < (u_0)^{(5)} < (\bar{u}_1)^{(5)}$,<br>and $\boxed{(u_0)^{(5)} = \frac{T_{28}^0}{T_{29}^0}}$<br><br>$(\mu_2)^{(5)} = (u_1)^{(5)}, (\mu_1)^{(5)} = (u_0)^{(5)}$, ***if*** $(\bar{u}_1)^{(5)} < (u_0)^{(5)}$ where $(u_1)^{(5)}, (\bar{u}_1)^{(5)}$ | 341 |
| Then the solution of global equations satisfies the inequalities<br><br>$G_{28}^0 e^{((S_1)^{(5)} - (p_{28})^{(5)})t} \le G_{28}(t) \le G_{28}^0 e^{(S_1)^{(5)}t}$<br><br>where $(p_i)^{(5)}$ is defined by equation | 342 |
| $\frac{1}{(m_5)^{(5)}} G_{28}^0 e^{((S_1)^{(5)} - (p_{28})^{(5)})t} \le G_{29}(t) \le \frac{1}{(m_2)^{(5)}} G_{28}^0 e^{(S_1)^{(5)}t}$ | 343 |
| $\left(\frac{(a_{30})^{(5)} G_{28}^0}{(m_1)^{(5)}((S_1)^{(5)} - (p_{28})^{(5)} - (S_2)^{(5)})}\left[e^{((S_1)^{(5)} - (p_{28})^{(5)})t} - e^{-(S_2)^{(5)}t}\right] + G_{30}^0 e^{-(S_2)^{(5)}t} \le G_{30}(t) \le$<br>*(a30)5G280(m2)5(S1)5−(a30′)5e(S1)5t−e−(a30′)5t+ G300e−(a30′)5t* | 344 |
| $\boxed{T_{28}^0 e^{(R_1)^{(5)}t} \le T_{28}(t) \le T_{28}^0 e^{((R_1)^{(5)} + (r_{28})^{(5)})t}}$ | 345 |

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| $\frac{1}{(\mu_1)^{(5)}}T_{28}^0 e^{(R_1)^{(5)}t} \le T_{28}(t) \le \frac{1}{(\mu_2)^{(5)}}T_{28}^0 e^{((R_1)^{(5)}+(r_{28})^{(5)})t}$ | 346 |
| $\frac{(b_{30})^{(5)}T_{28}^0}{(\mu_1)^{(5)}((R_1)^{(5)}-(b_{30}^{'})^{(5)})}\left[e^{(R_1)^{(5)}t} - e^{-(b_{30}^{'})^{(5)}t}\right] + T_{30}^0 e^{-(b_{30}^{'})^{(5)}t} \le T_{30}(t) \le$ $\frac{(a_{30})^{(5)}T_{28}^0}{(\mu_2)^{(5)}((R_1)^{(5)}+(r_{28})^{(5)}+(R_2)^{(5)})}\left[e^{((R_1)^{(5)}+(r_{28})^{(5)})t} - e^{-(R_2)^{(5)}t}\right] + T_{30}^0 e^{-(R_2)^{(5)}t}$ | 347 |
| **Definition of** $(S_1)^{(5)}, (S_2)^{(5)}, (R_1)^{(5)}, (R_2)^{(5)}$:- Where $(S_1)^{(5)} = (a_{28})^{(5)}(m_2)^{(5)} - (a_{28}^{'})^{(5)}$ $(S_2)^{(5)} = (a_{30})^{(5)} - (p_{30})^{(5)}$ $(R_1)^{(5)} = (b_{28})^{(5)}(\mu_2)^{(5)} - (b_{28}^{'})^{(5)}$ $(R_2)^{(5)} = (b_{30}^{'})^{(5)} - (r_{30})^{(5)}$ | 348 |
| **Behavior of the solutions of equation** **Theorem 2:** If we denote and define **Definition of** $(\sigma_1)^{(6)}, (\sigma_2)^{(6)}, (\tau_1)^{(6)}, (\tau_2)^{(6)}$ : $(\sigma_1)^{(6)}, (\sigma_2)^{(6)}, (\tau_1)^{(6)}, (\tau_2)^{(6)}$ four constants satisfying $-(\sigma_2)^{(6)} \le -(a_{32}^{'})^{(6)} + (a_{33}^{'})^{(6)} - (a_{32}^{''})^{(6)}(T_{33}, t) + (a_{33}^{''})^{(6)}(T_{33}, t) \le -(\sigma_1)^{(6)}$ $-(\tau_2)^{(6)} \le -(b_{32}^{'})^{(6)} + (b_{33}^{'})^{(6)} - (b_{32}^{''})^{(6)}\left((G_{35}), t\right) - (b_{33}^{''})^{(6)}\left((G_{35}), t\right) \le -(\tau_1)^{(6)}$ | 349 |
| **Definition of** $(v_1)^{(6)}, (v_2)^{(6)}, (u_1)^{(6)}, (u_2)^{(6)}, v^{(6)}, u^{(6)}$ : By $(v_1)^{(6)} > 0$, $(v_2)^{(6)} < 0$ and respectively $(u_1)^{(6)} > 0$, $(u_2)^{(6)} < 0$ the roots of the equations $(a_{33})^{(6)}\left(v^{(6)}\right)^2 + (\sigma_1)^{(6)}v^{(6)} - (a_{32})^{(6)} = 0$ and $(b_{33})^{(6)}\left(u^{(6)}\right)^2 + (\tau_1)^{(6)}u^{(6)} - (b_{32})^{(6)} = 0$ and | 350 |
| **Definition of** $(\bar{v}_1)^{(6)},, (\bar{v}_2)^{(6)}, (\bar{u}_1)^{(6)}, (\bar{u}_2)^{(6)}$ : By $(\bar{v}_1)^{(6)} > 0$, $(\bar{v}_2)^{(6)} < 0$ and respectively $(\bar{u}_1)^{(6)} > 0$, $(\bar{u}_2)^{(6)} < 0$ the roots of the equations $(a_{33})^{(6)}\left(v^{(6)}\right)^2 + (\sigma_2)^{(6)}v^{(6)} - (a_{32})^{(6)} = 0$ and $(b_{33})^{(6)}\left(u^{(6)}\right)^2 + (\tau_2)^{(6)}u^{(6)} - (b_{32})^{(6)} = 0$ **Definition of** $(m_1)^{(6)}, (m_2)^{(6)}, (\mu_1)^{(6)}, (\mu_2)^{(6)}, (v_0)^{(6)}$ :- If we define $(m_1)^{(6)}, (m_2)^{(6)}, (\mu_1)^{(6)}, (\mu_2)^{(6)}$ by $(m_2)^{(6)} = (v_0)^{(6)}, (m_1)^{(6)} = (v_1)^{(6)}$, ***if*** $(v_0)^{(6)} < (v_1)^{(6)}$ $(m_2)^{(6)} = (v_1)^{(6)}, (m_1)^{(6)} = (\bar{v}_6)^{(6)}$, ***if*** $(v_1)^{(6)} < (v_0)^{(6)} < (\bar{v}_1)^{(6)}$, | 351 |

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| and $\boxed{(v_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0}}$<br>$(m_2)^{(6)} = (v_1)^{(6)}, (m_1)^{(6)} = (v_0)^{(6)},\ \boldsymbol{if}\ (\bar{v}_1)^{(6)} < (v_0)^{(6)}$ | |
| and analogously<br>$(\mu_2)^{(6)} = (u_0)^{(6)}, (\mu_1)^{(6)} = (u_1)^{(6)},\ \boldsymbol{if}\ (u_0)^{(6)} < (u_1)^{(6)}$<br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (\bar{u}_1)^{(6)}, \boldsymbol{if}\ (u_1)^{(6)} < (u_0)^{(6)} < (\bar{u}_1)^{(6)},$<br>and $\boxed{(u_0)^{(6)} = \frac{T_{32}^0}{T_{33}^0}}$<br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (u_0)^{(6)}, \boldsymbol{if}\ (\bar{u}_1)^{(6)} < (u_0)^{(6)}$ where $(u_1)^{(6)}, (\bar{u}_1)^{(6)}$ | 352 |
| Then the solution of global equations satisfies the inequalities<br>$G_{32}^0 e^{((S_1)^{(6)} - (p_{32})^{(6)})t} \le G_{32}(t) \le G_{32}^0 e^{(S_1)^{(6)} t}$<br>where $(p_i)^{(6)}$ is defined by equation | 353 |
| $\frac{1}{(m_1)^{(6)}} G_{32}^0 e^{((S_1)^{(6)} - (p_{32})^{(6)})t} \le G_{33}(t) \le \frac{1}{(m_2)^{(6)}} G_{32}^0 e^{(S_1)^{(6)} t}$ | 354 |
| $\left(\frac{(a_{34})^{(6)} G_{32}^0}{(m_1)^{(6)}((S_1)^{(6)} - (p_{32})^{(6)} - (S_2)^{(6)})}\right)\left[e^{((S_1)^{(6)} - (p_{32})^{(6)})t} - e^{-(S_2)^{(6)} t}\right] + G_{34}^0 e^{-(S_2)^{(6)} t} \le G_{34}(t) \le$<br>*(a34)6G320(m2)6(S1)6−(a34′)6e(S1)6t−e−(a34′)6t+ G340e−(a34′)6t* | 355 |
| $\boxed{T_{32}^0 e^{(R_1)^{(6)} t} \le T_{32}(t) \le T_{32}^0 e^{((R_1)^{(6)} + (r_{32})^{(6)})t}}$ | 356 |
| $\frac{1}{(\mu_1)^{(6)}} T_{32}^0 e^{(R_1)^{(6)} t} \le T_{32}(t) \le \frac{1}{(\mu_2)^{(6)}} T_{32}^0 e^{((R_1)^{(6)} + (r_{32})^{(6)})t}$ | 357 |
| $\frac{(b_{34})^{(6)} T_{32}^0}{(\mu_1)^{(6)}((R_1)^{(6)} - (b'_{34})^{(6)})}\left[e^{(R_1)^{(6)} t} - e^{-(b'_{34})^{(6)} t}\right] + T_{34}^0 e^{-(b'_{34})^{(6)} t} \le T_{34}(t) \le$<br>$\frac{(a_{34})^{(6)} T_{32}^0}{(\mu_2)^{(6)}((R_1)^{(6)} + (r_{32})^{(6)} + (R_2)^{(6)})}\left[e^{((R_1)^{(6)} + (r_{32})^{(6)})t} - e^{-(R_2)^{(6)} t}\right] + T_{34}^0 e^{-(R_2)^{(6)} t}$ | 358 |
| **Definition of** $(S_1)^{(6)}, (S_2)^{(6)}, (R_1)^{(6)}, (R_2)^{(6)}$:-<br>Where $(S_1)^{(6)} = (a_{32})^{(6)}(m_2)^{(6)} - (a'_{32})^{(6)}$<br>$(S_2)^{(6)} = (a_{34})^{(6)} - (p_{34})^{(6)}$<br>$(R_1)^{(6)} = (b_{32})^{(6)}(\mu_2)^{(6)} - (b'_{32})^{(6)}$<br>$(R_2)^{(6)} = (b'_{34})^{(6)} - (r_{34})^{(6)}$ | 359 |
| **Behavior of the solutions of equation** | |

**Theorem 2:** If we denote and define

**Definition of** $(\sigma_1)^{(7)}, (\sigma_2)^{(7)}, (\tau_1)^{(7)}, (\tau_2)^{(7)}$ :

$(\sigma_1)^{(7)}, (\sigma_2)^{(7)}, (\tau_1)^{(7)}, (\tau_2)^{(7)}$ four constants satisfying

$$-(\sigma_2)^{(7)} \le -(a_{36}^{'})^{(7)} + (a_{37}^{'})^{(7)} - (a_{36}^{''})^{(7)}(T_{37}, t) + (a_{37}^{''})^{(7)}(T_{37}, t) \le -(\sigma_1)^{(7)}$$

$$-(\tau_2)^{(7)} \le -(b_{36}^{'})^{(7)} + (b_{37}^{'})^{(7)} - (b_{36}^{''})^{(7)}\big((G_{39}), t\big) - (b_{37}^{''})^{(7)}\big((G_{39}), t\big) \le -(\tau_1)^{(7)}$$

**Definition of** $(\nu_1)^{(7)}, (\nu_2)^{(7)}, (u_1)^{(7)}, (u_2)^{(7)}, \nu^{(7)}, u^{(7)}$ : 361

By $(\nu_1)^{(7)} > 0, (\nu_2)^{(7)} < 0$ and respectively $(u_1)^{(7)} > 0, (u_2)^{(7)} < 0$ the roots of the equations $(a_{37})^{(7)}\big(\nu^{(7)}\big)^2 + (\sigma_1)^{(7)}\nu^{(7)} - (a_{36})^{(7)} = 0$
and $(b_{37})^{(7)}\big(u^{(7)}\big)^2 + (\tau_1)^{(7)}u^{(7)} - (b_{36})^{(7)} = 0$ and

**Definition of** $(\bar{\nu}_1)^{(7)}, , (\bar{\nu}_2)^{(7)}, (\bar{u}_1)^{(7)}, (\bar{u}_2)^{(7)}$ : 362

By $(\bar{\nu}_1)^{(7)} > 0, (\bar{\nu}_2)^{(7)} < 0$ and respectively $(\bar{u}_1)^{(7)} > 0, (\bar{u}_2)^{(7)} < 0$ the roots of the equations $(a_{37})^{(7)}\big(\nu^{(7)}\big)^2 + (\sigma_2)^{(7)}\nu^{(7)} - (a_{36})^{(7)} = 0$
and $(b_{37})^{(7)}\big(u^{(7)}\big)^2 + (\tau_2)^{(7)}u^{(7)} - (b_{36})^{(7)} = 0$

**Definition of** $(m_1)^{(7)}, (m_2)^{(7)}, (\mu_1)^{(7)}, (\mu_2)^{(7)}, (\nu_0)^{(7)}$ :-

If we define $(m_1)^{(7)}, (m_2)^{(7)}, (\mu_1)^{(7)}, (\mu_2)^{(7)}$ by

$(m_2)^{(7)} = (\nu_0)^{(7)}, (m_1)^{(7)} = (\nu_1)^{(7)}$, ***if*** $(\nu_0)^{(7)} < (\nu_1)^{(7)}$

$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\bar{\nu}_1)^{(7)}$, ***if*** $(\nu_1)^{(7)} < (\nu_0)^{(7)} < (\bar{\nu}_1)^{(7)}$,
and $\boxed{(\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0}}$

$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\nu_0)^{(7)}$, ***if*** $(\bar{\nu}_1)^{(7)} < (\nu_0)^{(7)}$

and analogously 363

$(\mu_2)^{(7)} = (u_0)^{(7)}, (\mu_1)^{(7)} = (u_1)^{(7)}$, ***if*** $(u_0)^{(7)} < (u_1)^{(7)}$

$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (\bar{u}_1)^{(7)}$, ***if*** $(u_1)^{(7)} < (u_0)^{(7)} < (\bar{u}_1)^{(7)}$,
and $\boxed{(u_0)^{(7)} = \frac{T_{36}^0}{T_{37}^0}}$

$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (u_0)^{(7)}$, ***if*** $(\bar{u}_1)^{(7)} < (u_0)^{(7)}$ where $(u_1)^{(7)}, (\bar{u}_1)^{(7)}$

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| Then the solution of global equations satisfies the inequalities<br>$G_{36}^0 e^{((S_1)^{(7)}-(p_{36})^{(7)})t} \le G_{36}(t) \le G_{36}^0 e^{(S_1)^{(7)}t}$<br>where $(p_i)^{(7)}$ is defined by equation | 364 |
| $\frac{1}{(m_7)^{(7)}} G_{36}^0 e^{((S_1)^{(7)}-(p_{36})^{(7)})t} \le G_{37}(t) \le \frac{1}{(m_2)^{(7)}} G_{36}^0 e^{(S_1)^{(7)}t}$ | 365 |
| $(\frac{(a_{38})^{(7)} G_{36}^0}{(m_1)^{(7)}((S_1)^{(7)}-(p_{36})^{(7)}-(S_2)^{(7)})}\left[e^{((S_1)^{(7)}-(p_{36})^{(7)})t} - e^{-(S_2)^{(7)}t}\right] + G_{38}^0 e^{-(S_2)^{(7)}t} \le G_{38}(t) \le$<br>$\frac{(a_{38})^{(7)} G_{36}^0}{(m_2)^{(7)}((S_1)^{(7)}-(a_{38}')^{(7)})}[e^{(S_1)^{(7)}t} - e^{-(a_{38}')^{(7)}t}] + G_{38}^0 e^{-(a_{38}')^{(7)}t})$ | 366 |
| $\boxed{T_{36}^0 e^{(R_1)^{(7)}t} \le T_{36}(t) \le T_{36}^0 e^{((R_1)^{(7)}+(r_{36})^{(7)})t}}$ | 367 |
| $\frac{1}{(\mu_1)^{(7)}} T_{36}^0 e^{(R_1)^{(7)}t} \le T_{36}(t) \le \frac{1}{(\mu_2)^{(7)}} T_{36}^0 e^{((R_1)^{(7)}+(r_{36})^{(7)})t}$ | 368 |
| $\frac{(b_{38})^{(7)} T_{36}^0}{(\mu_1)^{(7)}((R_1)^{(7)}-(b_{38}')^{(7)})}\left[e^{(R_1)^{(7)}t} - e^{-(b_{38}')^{(7)}t}\right] + T_{38}^0 e^{-(b_{38}')^{(7)}t} \le T_{38}(t) \le$<br>$\frac{(a_{38})^{(7)} T_{36}^0}{(\mu_2)^{(7)}((R_1)^{(7)}+(r_{36})^{(7)}+(R_2)^{(7)})}\left[e^{((R_1)^{(7)}+(r_{36})^{(7)})t} - e^{-(R_2)^{(7)}t}\right] + T_{38}^0 e^{-(R_2)^{(7)}t}$ | 369 |
| **Definition of** $(S_1)^{(7)}, (S_2)^{(7)}, (R_1)^{(7)}, (R_2)^{(7)}$:-<br>Where $(S_1)^{(7)} = (a_{36})^{(7)}(m_2)^{(7)} - (a_{36}')^{(7)}$<br>$(S_2)^{(7)} = (a_{38})^{(7)} - (p_{38})^{(7)}$<br>$(R_1)^{(7)} = (b_{36})^{(7)}(\mu_2)^{(7)} - (b_{36}')^{(7)}$<br>$(R_2)^{(7)} = (b_{38}')^{(7)} - (r_{38})^{(7)}$ | 370 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ :<br>$(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ four constants satisfying<br>$-(\sigma_2)^{(8)} \le -(a_{40}')^{(8)} + (a_{41}')^{(8)} - (a_{40}'')^{(8)}(T_{41}, t) + (a_{41}'')^{(8)}(T_{41}, t) \le -(\sigma_1)^{(8)}$<br>$-(\tau_2)^{(8)} \le -(b_{40}')^{(8)} + (b_{41}')^{(8)} - (b_{40}'')^{(8)}((G_{43}), t) - (b_{41}'')^{(8)}((G_{43}), t) \le -(\tau_1)^{(8)}$ | 371 |

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| **Definition of** $(v_1)^{(8)}, (v_2)^{(8)}, (u_1)^{(8)}, (u_2)^{(8)}, v^{(8)}, u^{(8)}$ : <br><br> By $(v_1)^{(8)} > 0$ , $(v_2)^{(8)} < 0$ and respectively $(u_1)^{(8)} > 0$ , $(u_2)^{(8)} < 0$ the roots of the equations $(a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_1)^{(8)}v^{(8)} - (a_{40})^{(8)} = 0$ <br> and $(b_{41})^{(8)}\left(u^{(8)}\right)^2 + (\tau_1)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$ and | 372 |
| **Definition of** $(\bar{v}_1)^{(8)}$,, $(\bar{v}_2)^{(8)}, (\bar{u}_1)^{(8)}, (\bar{u}_2)^{(8)}$ : <br><br> By $(\bar{v}_1)^{(8)} > 0$ , $(\bar{v}_2)^{(8)} < 0$ and respectively $(\bar{u}_1)^{(8)} > 0$ , $(\bar{u}_2)^{(8)} < 0$ the <br><br> roots of the equations $(a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_2)^{(8)}v^{(8)} - (a_{40})^{(8)} = 0$ <br><br> and $(b_{41})^{(8)}\left(u^{(8)}\right)^2 + (\tau_2)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$ <br><br> **Definition of** $(m_1)^{(8)}$ , $(m_2)^{(8)}$ , $(\mu_1)^{(8)}, (\mu_2)^{(8)}, (v_0)^{(8)}$ :- <br><br> If we define $(m_1)^{(8)}$ , $(m_2)^{(8)}$ , $(\mu_1)^{(8)}, (\mu_2)^{(8)}$ by <br><br> $(m_2)^{(8)} = (v_0)^{(8)}, (m_1)^{(8)} = (v_1)^{(8)}$, ***if*** $(v_0)^{(8)} < (v_1)^{(8)}$ <br><br> $(m_2)^{(8)} = (v_1)^{(8)}, (m_1)^{(8)} = (\bar{v}_1)^{(8)}$ , ***if*** $(v_1)^{(8)} < (v_0)^{(8)} < (\bar{v}_1)^{(8)}$, <br> and $\boxed{(v_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}}$ <br><br> $(m_2)^{(8)} = (v_1)^{(8)}, (m_1)^{(8)} = (v_0)^{(8)}$, ***if*** $(\bar{v}_1)^{(8)} < (v_0)^{(8)}$ | |
| and analogously <br><br> $(\mu_2)^{(8)} = (u_0)^{(8)}, (\mu_1)^{(8)} = (u_1)^{(8)}$, ***if*** $(u_0)^{(8)} < (u_1)^{(8)}$ <br><br> $(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (\bar{u}_1)^{(8)}$ , ***if*** $(u_1)^{(8)} < (u_0)^{(8)} < (\bar{u}_1)^{(8)}$, <br> and $\boxed{(u_0)^{(8)} = \frac{T_{40}^0}{T_{41}^0}}$ <br><br> $(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (u_0)^{(8)}$, ***if*** $(\bar{u}_1)^{(8)} < (u_0)^{(8)}$ where $(u_1)^{(8)}, (\bar{u}_1)^{(8)}$ | 374 |
| Then the solution of global equations satisfies the inequalities <br><br> $G_{40}^0 e^{((S_1)^{(8)} - (p_{40})^{(8)})t} \leq G_{40}(t) \leq G_{40}^0 e^{(S_1)^{(8)}t}$ <br><br> where $(p_i)^{(8)}$ is defined by equation | 375 |

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| $\frac{1}{(m_1)^{(8)}} G_{40}^0 e^{((S_1)^{(8)}-(p_{40})^{(8)})t} \le G_{41}(t) \le \frac{1}{(m_2)^{(8)}} G_{40}^0 e^{(S_1)^{(8)}t}$ | 376 |
| $\left(\frac{(a_{42})^{(8)}G_{40}^0}{(m_1)^{(8)}((S_1)^{(8)}-(p_{40})^{(8)}-(S_2)^{(8)})}\left[e^{((S_1)^{(8)}-(p_{40})^{(8)})t} - e^{-(S_2)^{(8)}t}\right] + G_{42}^0 e^{-(S_2)^{(8)}t} \le G_{42}(t) \le \right.$ $\left.\frac{(a_{42})^{(8)}G_{40}^0}{(m_2)^{(8)}((S_1)^{(8)}-(a'_{42})^{(8)})}[e^{(S_1)^{(8)}t} - e^{-(a'_{42})^{(8)}t}] + G_{42}^0 e^{-(a'_{42})^{(8)}t}\right)$ | 377 |
| $\boxed{T_{40}^0 e^{(R_1)^{(8)}t} \le T_{40}(t) \le T_{40}^0 e^{((R_1)^{(8)}+(r_{40})^{(8)})t}}$ | 378 |
| $\frac{1}{(\mu_1)^{(8)}} T_{40}^0 e^{(R_1)^{(8)}t} \le T_{40}(t) \le \frac{1}{(\mu_2)^{(8)}} T_{40}^0 e^{((R_1)^{(8)}+(r_{40})^{(8)})t}$ | 379 |
| $\frac{(b_{42})^{(8)}T_{40}^0}{(\mu_1)^{(8)}((R_1)^{(8)}-(b'_{42})^{(8)})}\left[e^{(R_1)^{(8)}t} - e^{-(b'_{42})^{(8)}t}\right] + T_{42}^0 e^{-(b'_{42})^{(8)}t} \le T_{42}(t) \le$ $\frac{(a_{42})^{(8)}T_{40}^0}{(\mu_2)^{(8)}((R_1)^{(8)}+(r_{40})^{(8)}+(R_2)^{(8)})}\left[e^{((R_1)^{(8)}+(r_{40})^{(8)})t} - e^{-(R_2)^{(8)}t}\right] + T_{42}^0 e^{-(R_2)^{(8)}t}$ | 380 |
| **Definition of** $(S_1)^{(8)}, (S_2)^{(8)}, (R_1)^{(8)}, (R_2)^{(8)}$:-<br>Where $(S_1)^{(8)} = (a_{40})^{(8)}(m_2)^{(8)} - (a'_{40})^{(8)}$<br>$(S_2)^{(8)} = (a_{42})^{(8)} - (p_{42})^{(8)}$<br>$(R_1)^{(8)} = (b_{40})^{(8)}(\mu_2)^{(8)} - (b'_{40})^{(8)}$<br>$(R_2)^{(8)} = (b'_{42})^{(8)} - (r_{42})^{(8)}$ | 381 |
| **Behavior of the solutions of equation 37 to 92**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(9)}, (\sigma_2)^{(9)}, (\tau_1)^{(9)}, (\tau_2)^{(9)}$ :<br>$(\sigma_1)^{(9)}, (\sigma_2)^{(9)}, (\tau_1)^{(9)}, (\tau_2)^{(9)}$ four constants satisfying<br>$-(\sigma_2)^{(9)} \le -(a'_{44})^{(9)} + (a'_{45})^{(9)} - (a''_{44})^{(9)}(T_{45}, t) + (a''_{45})^{(9)}(T_{45}, t) \le -(\sigma_1)^{(9)}$<br>$-(\tau_2)^{(9)} \le -(b'_{44})^{(9)} + (b'_{45})^{(9)} - (b''_{44})^{(9)}((G_{47}), t) - (b''_{45})^{(9)}((G_{47}), t) \le -(\tau_1)^{(9)}$ | 382 |
| **Definition of** $(\nu_1)^{(9)}, (\nu_2)^{(9)}, (u_1)^{(9)}, (u_2)^{(9)}, \nu^{(9)}, u^{(9)}$ :<br>By $(\nu_1)^{(9)} > 0$, $(\nu_2)^{(9)} < 0$ and respectively $(u_1)^{(9)} > 0$, $(u_2)^{(9)} < 0$ the roots of the equations $(a_{45})^{(9)}(\nu^{(9)})^2 + (\sigma_1)^{(9)}\nu^{(9)} - (a_{44})^{(9)} = 0$<br>and $(b_{45})^{(9)}(u^{(9)})^2 + (\tau_1)^{(9)}u^{(9)} - (b_{44})^{(9)} = 0$ and | |
| **Definition of** $(\bar{\nu}_1)^{(9)},, (\bar{\nu}_2)^{(9)}, (\bar{u}_1)^{(9)}, (\bar{u}_2)^{(9)}$ : | |

By $(\bar{v}_1)^{(9)} > 0$ , $(\bar{v}_2)^{(9)} < 0$ and respectively $(\bar{u}_1)^{(9)} > 0$ , $(\bar{u}_2)^{(9)} < 0$ the roots of the equations $(a_{45})^{(9)}\left(v^{(9)}\right)^2 + (\sigma_2)^{(9)}v^{(9)} - (a_{44})^{(9)} = 0$ and $(b_{45})^{(9)}\left(u^{(9)}\right)^2 + (\tau_2)^{(9)}u^{(9)} - (b_{44})^{(9)} = 0$

**Definition of** $(m_1)^{(9)}$ , $(m_2)^{(9)}$ , $(\mu_1)^{(9)}$, $(\mu_2)^{(9)}$, $(v_0)^{(9)}$ :-

If we define $(m_1)^{(9)}$ , $(m_2)^{(9)}$ , $(\mu_1)^{(9)}$, $(\mu_2)^{(9)}$ by

$(m_2)^{(9)} = (v_0)^{(9)}, (m_1)^{(9)} = (v_1)^{(9)}$, ***if*** $(v_0)^{(9)} < (v_1)^{(9)}$

$(m_2)^{(9)} = (v_1)^{(9)}, (m_1)^{(9)} = (\bar{v}_1)^{(9)}$ , ***if*** $(v_1)^{(9)} < (v_0)^{(9)} < (\bar{v}_1)^{(9)}$,

and $\boxed{(v_0)^{(9)} = \frac{G_{44}^0}{G_{45}^0}}$

$(m_2)^{(9)} = (v_1)^{(9)}, (m_1)^{(9)} = (v_0)^{(9)}$, ***if*** $(\bar{v}_1)^{(9)} < (v_0)^{(9)}$

and analogously

$(\mu_2)^{(9)} = (u_0)^{(9)}, (\mu_1)^{(9)} = (u_1)^{(9)}$, ***if*** $(u_0)^{(9)} < (u_1)^{(9)}$

$(\mu_2)^{(9)} = (u_1)^{(9)}, (\mu_1)^{(9)} = (\bar{u}_1)^{(9)}$ , ***if*** $(u_1)^{(9)} < (u_0)^{(9)} < (\bar{u}_1)^{(9)}$,

and $\boxed{(u_0)^{(9)} = \frac{T_{44}^0}{T_{45}^0}}$

$(\mu_2)^{(9)} = (u_1)^{(9)}, (\mu_1)^{(9)} = (u_0)^{(9)}$, ***if*** $(\bar{u}_1)^{(9)} < (u_0)^{(9)}$ where $(u_1)^{(9)}, (\bar{u}_1)^{(9)}$ are defined by 59 and 69 respectively

Then the solution of 19,20,21,22,23 and 24 satisfies the inequalities

$$G_{44}^0 e^{((S_1)^{(9)} - (p_{44})^{(9)})t} \le G_{44}(t) \le G_{44}^0 e^{(S_1)^{(9)}t}$$

where $(p_i)^{(9)}$ is defined by equation 45

$$\frac{1}{(m_9)^{(9)}} G_{44}^0 e^{((S_1)^{(9)} - (p_{44})^{(9)})t} \le G_{45}(t) \le \frac{1}{(m_2)^{(9)}} G_{44}^0 e^{(S_1)^{(9)}t}$$

$$\left( \frac{(a_{46})^{(9)} G_{44}^0}{(m_1)^{(9)}((S_1)^{(9)} - (p_{44})^{(9)} - (S_2)^{(9)})} \left[ e^{((S_1)^{(9)} - (p_{44})^{(9)})t} - e^{-(S_2)^{(9)}t} \right] + G_{46}^0 e^{-(S_2)^{(9)}t} \le G_{46}(t) \le \right.$$

$$\left. \frac{(a_{46})^{(9)} G_{44}^0}{(m_2)^{(9)}((S_1)^{(9)} - (a'_{46})^{(9)})} [e^{(S_1)^{(9)}t} - e^{-(a'_{46})^{(9)}t}] + G_{46}^0 e^{-(a'_{46})^{(9)}t} \right)$$

$$\boxed{T_{44}^0 e^{(R_1)^{(9)}t} \le T_{44}(t) \le T_{44}^0 e^{((R_1)^{(9)} + (r_{44})^{(9)})t}}$$

$$\frac{1}{(\mu_1)^{(9)}} T_{44}^0 e^{(R_1)^{(9)}t} \le T_{44}(t) \le \frac{1}{(\mu_2)^{(9)}} T_{44}^0 e^{((R_1)^{(9)} + (r_{44})^{(9)})t}$$

$$\frac{(b_{46})^{(9)} T_{44}^0}{(\mu_1)^{(9)}((R_1)^{(9)} - (b'_{46})^{(9)})} \left[ e^{(R_1)^{(9)}t} - e^{-(b'_{46})^{(9)}t} \right] + T_{46}^0 e^{-(b'_{46})^{(9)}t} \le T_{46}(t) \le$$

$\frac{(a_{46})^{(9)}T_{44}^{0}}{(\mu_2)^{(9)}((R_1)^{(9)}+(r_{44})^{(9)}+(R_2)^{(9)})}\left[e^{((R_1)^{(9)}+(r_{44})^{(9)})t}-e^{-(R_2)^{(9)}t}\right]+T_{46}^{0}e^{-(R_2)^{(9)}t}$

**Definition of** $(S_1)^{(9)},(S_2)^{(9)},(R_1)^{(9)},(R_2)^{(9)}$:-

Where $(S_1)^{(9)}=(a_{44})^{(9)}(m_2)^{(9)}-(a_{44}^{'})^{(9)}$

$(S_2)^{(9)}=(a_{46})^{(9)}-(p_{46})^{(9)}$

$(R_1)^{(9)}=(b_{44})^{(9)}(\mu_2)^{(9)}-(b_{44}^{'})^{(9)}$

$(R_2)^{(9)}=(b_{46}^{'})^{(9)}-(r_{46})^{(9)}$

**Proof :** From global equations we obtain

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$$\frac{dv^{(1)}}{dt}=(a_{13})^{(1)}-\left((a_{13}^{'})^{(1)}-(a_{14}^{'})^{(1)}+(a_{13}^{''})^{(1)}(T_{14},t)\right)-(a_{14}^{''})^{(1)}(T_{14},t)v^{(1)}-(a_{14})^{(1)}v^{(1)}$$

**Definition of** $v^{(1)}$ :- $\boxed{v^{(1)}=\frac{G_{13}}{G_{14}}}$

It follows

$$-\left((a_{14})^{(1)}\left(v^{(1)}\right)^2+(\sigma_2)^{(1)}v^{(1)}-(a_{13})^{(1)}\right)\leq\frac{dv^{(1)}}{dt}\leq-\left((a_{14})^{(1)}\left(v^{(1)}\right)^2+(\sigma_1)^{(1)}v^{(1)}-(a_{13})^{(1)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(1)},(\nu_0)^{(1)}$ :-

For $0<\boxed{(\nu_0)^{(1)}=\frac{G_{13}^{0}}{G_{14}^{0}}}<(\nu_1)^{(1)}<(\bar{\nu}_1)^{(1)}$

$$\nu^{(1)}(t)\geq\frac{(\nu_1)^{(1)}+(C)^{(1)}(\nu_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)}-(\nu_0)^{(1)}\right)t\right]}}{1+(C)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)}-(\nu_0)^{(1)}\right)t\right]}}\quad,\quad\boxed{(C)^{(1)}=\frac{(\nu_1)^{(1)}-(\nu_0)^{(1)}}{(\nu_0)^{(1)}-(\nu_2)^{(1)}}}$$

it follows $(\nu_0)^{(1)}\leq\nu^{(1)}(t)\leq(\nu_1)^{(1)}$

In the same manner , we get

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$$\nu^{(1)}(t)\leq\frac{(\bar{\nu}_1)^{(1)}+(\bar{C})^{(1)}(\bar{\nu}_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}}{1+(\bar{C})^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}}\quad,\quad\boxed{(\bar{C})^{(1)}=\frac{(\bar{\nu}_1)^{(1)}-(\nu_0)^{(1)}}{(\nu_0)^{(1)}-(\bar{\nu}_2)^{(1)}}}$$

From which we deduce $(\nu_0)^{(1)}\leq\nu^{(1)}(t)\leq(\bar{\nu}_1)^{(1)}$

If $0<(\nu_1)^{(1)}<(\nu_0)^{(1)}=\frac{G_{13}^{0}}{G_{14}^{0}}<(\bar{\nu}_1)^{(1)}$ we find like in the previous case,

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$$(\nu_1)^{(1)} \le \frac{(\nu_1)^{(1)}+(C)^{(1)}(\nu_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)}-(\nu_2)^{(1)}\right)t\right]}}{1+(C)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)}-(\nu_2)^{(1)}\right)t\right]}} \le \nu^{(1)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(1)}+(\bar{C})^{(1)}(\bar{\nu}_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}}{1+(\bar{C})^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}} \le (\bar{\nu}_1)^{(1)}$$

If $0 < (\nu_1)^{(1)} \le (\bar{\nu}_1)^{(1)} \le \boxed{(\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}}$ , we obtain 386

$$(\nu_1)^{(1)} \le \nu^{(1)}(t) \le \frac{(\bar{\nu}_1)^{(1)}+(\bar{C})^{(1)}(\bar{\nu}_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}}{1+(\bar{C})^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}} \le (\nu_0)^{(1)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(1)}(t)$ :-

$$(m_2)^{(1)} \le \nu^{(1)}(t) \le (m_1)^{(1)}, \quad \boxed{\nu^{(1)}(t) = \frac{G_{13}(t)}{G_{14}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(1)}(t)$ :-

$$(\mu_2)^{(1)} \le u^{(1)}(t) \le (\mu_1)^{(1)}, \quad \boxed{u^{(1)}(t) = \frac{T_{13}(t)}{T_{14}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{13}^{''})^{(1)} = (a_{14}^{''})^{(1)}, then\ (\sigma_1)^{(1)} = (\sigma_2)^{(1)}$ and in this case $(\nu_1)^{(1)} = (\bar{\nu}_1)^{(1)}$ if in addition $(\nu_0)^{(1)} = (\nu_1)^{(1)}$ then $\nu^{(1)}(t) = (\nu_0)^{(1)}$ and as a consequence $G_{13}(t) = (\nu_0)^{(1)}G_{14}(t)$ this also defines $(\nu_0)^{(1)}$ for the special case

Analogously if $(b_{13}^{''})^{(1)} = (b_{14}^{''})^{(1)}, then\ (\tau_1)^{(1)} = (\tau_2)^{(1)}$ and then

$(u_1)^{(1)} = (\bar{u}_1)^{(1)}$if in addition $(u_0)^{(1)} = (u_1)^{(1)}$ then $T_{13}(t) = (u_0)^{(1)}T_{14}(t)$ This is an important consequence of the relation between $(\nu_1)^{(1)}$ and $(\bar{\nu}_1)^{(1)}$, and definition of $(u_0)^{(1)}$.

**Proof** : From global equations we obtain 387

$$\frac{d\nu^{(2)}}{dt} = (a_{16})^{(2)} - \left((a_{16}')^{(2)} - (a_{17}')^{(2)} + (a_{16}'')^{(2)}(T_{17}, t)\right) - (a_{17}'')^{(2)}(T_{17}, t)\nu^{(2)} - (a_{17})^{(2)}\nu^{(2)}$$

**Definition of** $\nu^{(2)}$ :- $\boxed{\nu^{(2)} = \frac{G_{16}}{G_{17}}}$ 388

It follows 389

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| $-\left((a_{17})^{(2)}\left(v^{(2)}\right)^2+(\sigma_2)^{(2)}v^{(2)}-(a_{16})^{(2)}\right)\le\frac{dv^{(2)}}{dt}\le-\left((a_{17})^{(2)}\left(v^{(2)}\right)^2+(\sigma_1)^{(2)}v^{(2)}-(a_{16})^{(2)}\right)$ | |
| From which one obtains<br><br>**Definition of** $(\bar{v}_1)^{(2)},(v_0)^{(2)}$ :-<br><br>For $0<(v_0)^{(2)}=\frac{G_{16}^0}{G_{17}^0}<(v_1)^{(2)}<(\bar{v}_1)^{(2)}$<br><br>$v^{(2)}(t)\ge\frac{(v_1)^{(2)}+(C)^{(2)}(v_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((v_1)^{(2)}-(v_0)^{(2)}\right)t\right]}}{1+(C)^{(2)}e^{\left[-(a_{17})^{(2)}\left((v_1)^{(2)}-(v_0)^{(2)}\right)t\right]}}$ , $\boxed{(C)^{(2)}=\frac{(v_1)^{(2)}-(v_0)^{(2)}}{(v_0)^{(2)}-(v_2)^{(2)}}}$<br><br>it follows $(v_0)^{(2)}\le v^{(2)}(t)\le(v_1)^{(2)}$ | 390 |
| In the same manner , we get<br><br>$v^{(2)}(t)\le\frac{(\bar{v}_1)^{(2)}+(\bar{C})^{(2)}(\bar{v}_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}{1+(\bar{C})^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}$ , $\boxed{(\bar{C})^{(2)}=\frac{(\bar{v}_1)^{(2)}-(v_0)^{(2)}}{(v_0)^{(2)}-(\bar{v}_2)^{(2)}}}$ | 391 |
| From which we deduce $(v_0)^{(2)}\le v^{(2)}(t)\le(\bar{v}_1)^{(2)}$ | 392 |
| If $0<(v_1)^{(2)}<(v_0)^{(2)}=\frac{G_{16}^0}{G_{17}^0}<(\bar{v}_1)^{(2)}$ we find like in the previous case,<br><br>$(v_1)^{(2)}\le\frac{(v_1)^{(2)}+(C)^{(2)}(v_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((v_1)^{(2)}-(v_2)^{(2)}\right)t\right]}}{1+(C)^{(2)}e^{\left[-(a_{17})^{(2)}\left((v_1)^{(2)}-(v_2)^{(2)}\right)t\right]}}\le v^{(2)}(t)\le$<br><br>$\frac{(\bar{v}_1)^{(2)}+(\bar{C})^{(2)}(\bar{v}_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}{1+(\bar{C})^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}\le(\bar{v}_1)^{(2)}$ | 393 |
| If $0<(v_1)^{(2)}\le(\bar{v}_1)^{(2)}\le(v_0)^{(2)}=\frac{G_{16}^0}{G_{17}^0}$ , we obtain<br><br>$(v_1)^{(2)}\le v^{(2)}(t)\le\frac{(\bar{v}_1)^{(2)}+(\bar{C})^{(2)}(\bar{v}_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}{1+(\bar{C})^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}\le(v_0)^{(2)}$<br><br>And so with the notation of the first part of condition (c) , we have | 394 |
| **Definition of** $v^{(2)}(t)$ :-<br><br>$(m_2)^{(2)}\le v^{(2)}(t)\le(m_1)^{(2)}$, $\boxed{v^{(2)}(t)=\frac{G_{16}(t)}{G_{17}(t)}}$ | 395 |
| In a completely analogous way, we obtain<br><br>**Definition of** $u^{(2)}(t)$ :- | 396 |

| | |
|---|---|
| $(\mu_2)^{(2)} \leq u^{(2)}(t) \leq (\mu_1)^{(2)}$, $\boxed{u^{(2)}(t) = \frac{T_{16}(t)}{T_{17}(t)}}$ | |
| Now, using this result and replacing it in global equations we get easily the result stated in the theorem. | |
| **Particular case :**<br><br>If $(a_{16}^{''})^{(2)} = (a_{17}^{''})^{(2)}, then\ (\sigma_1)^{(2)} = (\sigma_2)^{(2)}$ and in this case $(\nu_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ if in addition $(\nu_0)^{(2)} = (\nu_1)^{(2)}$ then $\nu^{(2)}(t) = (\nu_0)^{(2)}$ and as a consequence $G_{16}(t) = (\nu_0)^{(2)} G_{17}(t)$<br><br>Analogously if $(b_{16}^{''})^{(2)} = (b_{17}^{''})^{(2)}, then\ (\tau_1)^{(2)} = (\tau_2)^{(2)}$ and then<br><br>$(u_1)^{(2)} = (\bar{u}_1)^{(2)}$if in addition $(u_0)^{(2)} = (u_1)^{(2)}$ then $T_{16}(t) = (u_0)^{(2)} T_{17}(t)$ This is an important consequence of the relation between $(\nu_1)^{(2)}$ and $(\bar{\nu}_1)^{(2)}$ | 397 |
| **Proof** : From global equations we obtain<br><br>$\frac{d\nu^{(3)}}{dt} = (a_{20})^{(3)} - \left((a_{20}^{'})^{(3)} - (a_{21}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21}, t)\right) - (a_{21}^{''})^{(3)}(T_{21}, t)\nu^{(3)} - (a_{21})^{(3)}\nu^{(3)}$ | 398 |
| **Definition of** $\nu^{(3)}$ :- $\boxed{\nu^{(3)} = \frac{G_{20}}{G_{21}}}$<br><br>It follows<br><br>$-\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_2)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right) \leq \frac{d\nu^{(3)}}{dt} \leq -\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_1)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right)$ | 399 |
| From which one obtains<br><br>For $0 < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\nu_1)^{(3)} < (\bar{\nu}_1)^{(3)}$<br><br>$\nu^{(3)}(t) \geq \frac{(\nu_1)^{(3)} + (C)^{(3)}(\nu_2)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)} - (\nu_0)^{(3)}\right)t\right]}}{1 + (C)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)} - (\nu_0)^{(3)}\right)t\right]}}$ , $\boxed{(C)^{(3)} = \frac{(\nu_1)^{(3)} - (\nu_0)^{(3)}}{(\nu_0)^{(3)} - (\nu_2)^{(3)}}}$<br><br>it follows $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\nu_1)^{(3)}$ | 400 |
| In the same manner , we get<br><br>$\nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)} + (C)^{(3)}(\bar{\nu}_2)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)} - (\bar{\nu}_2)^{(3)}\right)t\right]}}{1 + (C)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)} - (\bar{\nu}_2)^{(3)}\right)t\right]}}$ , $\boxed{(\bar{C})^{(3)} = \frac{(\bar{\nu}_1)^{(3)} - (\nu_0)^{(3)}}{(\nu_0)^{(3)} - (\bar{\nu}_2)^{(3)}}}$<br><br>**Definition of** $(\bar{\nu}_1)^{(3)}$ :-<br><br>From which we deduce $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\bar{\nu}_1)^{(3)}$ | 401 |

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If $0 < (\nu_1)^{(3)} < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\bar{\nu}_1)^{(3)}$ we find like in the previous case,

$$(\nu_1)^{(3)} \leq \frac{(\nu_1)^{(3)}+(C)^{(3)}(\nu_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}}{1+(C)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}} \leq \nu^{(3)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(C)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\bar{\nu}_1)^{(3)}$$

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If $0 < (\nu_1)^{(3)} \leq (\bar{\nu}_1)^{(3)} \leq (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0}$ , we obtain

$$(\nu_1)^{(3)} \leq \nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\nu_0)^{(3)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(3)}(t)$ :-

$$(m_2)^{(3)} \leq \nu^{(3)}(t) \leq (m_1)^{(3)}, \quad \boxed{\nu^{(3)}(t) = \frac{G_{20}(t)}{G_{21}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(3)}(t)$ :-

$$(\mu_2)^{(3)} \leq u^{(3)}(t) \leq (\mu_1)^{(3)}, \quad \boxed{u^{(3)}(t) = \frac{T_{20}(t)}{T_{21}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{20}^{''})^{(3)} = (a_{21}^{''})^{(3)}, then\ (\sigma_1)^{(3)} = (\sigma_2)^{(3)}$ and in this case $(\nu_1)^{(3)} = (\bar{\nu}_1)^{(3)}$ if in addition $(\nu_0)^{(3)} = (\nu_1)^{(3)}$ then $\nu^{(3)}(t) = (\nu_0)^{(3)}$ and as a consequence $G_{20}(t) = (\nu_0)^{(3)}G_{21}(t)$

Analogously if $(b_{20}^{''})^{(3)} = (b_{21}^{''})^{(3)}, then\ (\tau_1)^{(3)} = (\tau_2)^{(3)}$ and then

$(u_1)^{(3)} = (\bar{u}_1)^{(3)}$if in addition $(u_0)^{(3)} = (u_1)^{(3)}$ then $T_{20}(t) = (u_0)^{(3)}T_{21}(t)$ This is an important consequence of the relation between $(\nu_1)^{(3)}$ and $(\bar{\nu}_1)^{(3)}$

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(4)}}{dt} = (a_{24})^{(4)} - \left((a_{24}^{'})^{(4)} - (a_{25}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25},t)\right) - (a_{25}^{''})^{(4)}(T_{25},t)\nu^{(4)} - (a_{25})^{(4)}\nu^{(4)}$$

**Definition of** $\nu^{(4)}$ :- $\boxed{\nu^{(4)} = \frac{G_{24}}{G_{25}}}$

It follows

$$-\left((a_{25})^{(4)}\left(v^{(4)}\right)^2+(\sigma_2)^{(4)}v^{(4)}-(a_{24})^{(4)}\right)\le\frac{dv^{(4)}}{dt}\le-\left((a_{25})^{(4)}\left(v^{(4)}\right)^2+(\sigma_4)^{(4)}v^{(4)}-(a_{24})^{(4)}\right)$$

From which one obtains

**Definition of** $(\bar{v}_1)^{(4)},(v_0)^{(4)}$ :-

For $0<\boxed{(v_0)^{(4)}=\frac{G_{24}^0}{G_{25}^0}}<(v_1)^{(4)}<(\bar{v}_1)^{(4)}$

$$v^{(4)}(t)\ge\frac{(v_1)^{(4)}+(C)^{(4)}(v_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((v_1)^{(4)}-(v_0)^{(4)}\right)t\right]}}{4+(C)^{(4)}e^{\left[-(a_{25})^{(4)}\left((v_1)^{(4)}-(v_0)^{(4)}\right)t\right]}}\quad,\quad\boxed{(C)^{(4)}=\frac{(v_1)^{(4)}-(v_0)^{(4)}}{(v_0)^{(4)}-(v_2)^{(4)}}}$$

it follows $(v_0)^{(4)}\le v^{(4)}(t)\le(v_1)^{(4)}$

In the same manner , we get

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$$v^{(4)}(t)\le\frac{(\bar{v}_1)^{(4)}+(\bar{C})^{(4)}(\bar{v}_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}{4+(\bar{C})^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}\quad,\quad\boxed{(\bar{C})^{(4)}=\frac{(\bar{v}_1)^{(4)}-(v_0)^{(4)}}{(v_0)^{(4)}-(\bar{v}_2)^{(4)}}}$$

From which we deduce $(v_0)^{(4)}\le v^{(4)}(t)\le(\bar{v}_1)^{(4)}$

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If $0<(v_1)^{(4)}<(v_0)^{(4)}=\frac{G_{24}^0}{G_{25}^0}<(\bar{v}_1)^{(4)}$ we find like in the previous case,

$$(v_1)^{(4)}\le\frac{(v_1)^{(4)}+(C)^{(4)}(v_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((v_1)^{(4)}-(v_2)^{(4)}\right)t\right]}}{1+(C)^{(4)}e^{\left[-(a_{25})^{(4)}\left((v_1)^{(4)}-(v_2)^{(4)}\right)t\right]}}\le v^{(4)}(t)\le$$

$$\frac{(\bar{v}_1)^{(4)}+(\bar{C})^{(4)}(\bar{v}_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}{1+(\bar{C})^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}\le(\bar{v}_1)^{(4)}$$

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If $0<(v_1)^{(4)}\le(\bar{v}_1)^{(4)}\le\boxed{(v_0)^{(4)}=\frac{G_{24}^0}{G_{25}^0}}$ , we obtain

$$(v_1)^{(4)}\le v^{(4)}(t)\le\frac{(\bar{v}_1)^{(4)}+(\bar{C})^{(4)}(\bar{v}_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}{1+(\bar{C})^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}\le(v_0)^{(4)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $v^{(4)}(t)$ :-

$$(m_2)^{(4)}\le v^{(4)}(t)\le(m_1)^{(4)},\quad\boxed{v^{(4)}(t)=\frac{G_{24}(t)}{G_{25}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(4)}(t)$ :-

$$(\mu_2)^{(4)}\le u^{(4)}(t)\le(\mu_1)^{(4)},\quad\boxed{u^{(4)}(t)=\frac{T_{24}(t)}{T_{25}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{24}^{''})^{(4)} = (a_{25}^{''})^{(4)}, then\ (\sigma_1)^{(4)} = (\sigma_2)^{(4)}$ and in this case $(\nu_1)^{(4)} = (\bar{\nu}_1)^{(4)}$ if in addition $(\nu_0)^{(4)} = (\nu_1)^{(4)}$ then $\nu^{(4)}(t) = (\nu_0)^{(4)}$ and as a consequence $G_{24}(t) = (\nu_0)^{(4)} G_{25}(t)$ **this also defines** $(\nu_0)^{(4)}$ **for the special case** .

Analogously if $(b_{24}^{''})^{(4)} = (b_{25}^{''})^{(4)}, then\ (\tau_1)^{(4)} = (\tau_2)^{(4)}$ and then $(u_1)^{(4)} = (\bar{u}_4)^{(4)}$ if in addition $(u_0)^{(4)} = (u_1)^{(4)}$ then $T_{24}(t) = (u_0)^{(4)} T_{25}(t)$ This is an important consequence of the relation between $(\nu_1)^{(4)}$ and $(\bar{\nu}_1)^{(4)}$, **and definition of** $(u_0)^{(4)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(5)}}{dt} = (a_{28})^{(5)} - \left((a_{28}^{'})^{(5)} - (a_{29}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}, t)\right) - (a_{29}^{''})^{(5)}(T_{29}, t)\nu^{(5)} - (a_{29})^{(5)}\nu^{(5)}$$

**Definition of** $\nu^{(5)}$ :- $\boxed{\nu^{(5)} = \frac{G_{28}}{G_{29}}}$

It follows

$$-\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_2)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right) \leq \frac{d\nu^{(5)}}{dt} \leq -\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_1)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(5)}, (\nu_0)^{(5)}$ :-

For $0 < \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}} < (\nu_1)^{(5)} < (\bar{\nu}_1)^{(5)}$

$$\nu^{(5)}(t) \geq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}}{5 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}} \quad , \quad \boxed{(C)^{(5)} = \frac{(\nu_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\nu_2)^{(5)}}}$$

it follows $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\nu_1)^{(5)}$

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In the same manner , we get

$$\nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)} + (\bar{C})^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{5 + (\bar{C})^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(5)} = \frac{(\bar{\nu}_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\bar{\nu}_2)^{(5)}}}$$

From which we deduce $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\bar{\nu}_5)^{(5)}$

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If $0 < (\nu_1)^{(5)} < (\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0} < (\bar{\nu}_1)^{(5)}$ we find like in the previous case,

$$(\nu_1)^{(5)} \leq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}}{1 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}} \leq \nu^{(5)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(5)}+(\bar{C})^{(5)}(\bar{\nu}_2)^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}}{1+(\bar{C})^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\bar{\nu}_1)^{(5)}$$

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If $0 < (\nu_1)^{(5)} \leq (\bar{\nu}_1)^{(5)} \leq \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}}$ , we obtain

$$(\nu_1)^{(5)} \leq \nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)}+(\bar{C})^{(5)}(\bar{\nu}_2)^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}}{1+(\bar{C})^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\nu_0)^{(5)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(5)}(t)$ :-

$(m_2)^{(5)} \leq \nu^{(5)}(t) \leq (m_1)^{(5)}$, $\boxed{\nu^{(5)}(t) = \frac{G_{28}(t)}{G_{29}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(5)}(t)$ :-

$(\mu_2)^{(5)} \leq u^{(5)}(t) \leq (\mu_1)^{(5)}$, $\boxed{u^{(5)}(t) = \frac{T_{28}(t)}{T_{29}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{28}^{''})^{(5)} = (a_{29}^{''})^{(5)}, then\ (\sigma_1)^{(5)} = (\sigma_2)^{(5)}$ and in this case $(\nu_1)^{(5)} = (\bar{\nu}_1)^{(5)}$ if in addition $(\nu_0)^{(5)} = (\nu_5)^{(5)}$ then $\nu^{(5)}(t) = (\nu_0)^{(5)}$ and as a consequence $G_{28}(t) = (\nu_0)^{(5)}G_{29}(t)$ **this also defines** $(\nu_0)^{(5)}$ **for the special case** .

Analogously if $(b_{28}^{''})^{(5)} = (b_{29}^{''})^{(5)}, then\ (\tau_1)^{(5)} = (\tau_2)^{(5)}$ and then $(u_1)^{(5)} = (\bar{u}_1)^{(5)}$ if in addition $(u_0)^{(5)} = (u_1)^{(5)}$ then $T_{28}(t) = (u_0)^{(5)}T_{29}(t)$ This is an important consequence of the relation between $(\nu_1)^{(5)}$ and $(\bar{\nu}_1)^{(5)}$, **and definition of** $(u_0)^{(5)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(6)}}{dt} = (a_{32})^{(6)} - \left((a_{32}^{'})^{(6)} - (a_{33}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33},t)\right) - (a_{33}^{''})^{(6)}(T_{33},t)\nu^{(6)} - (a_{33})^{(6)}\nu^{(6)}$$

**Definition of** $\nu^{(6)}$ :- $\boxed{\nu^{(6)} = \frac{G_{32}}{G_{33}}}$

It follows

$$-\left((a_{33})^{(6)}\left(\nu^{(6)}\right)^2 + (\sigma_2)^{(6)}\nu^{(6)} - (a_{32})^{(6)}\right) \leq \frac{d\nu^{(6)}}{dt} \leq -\left((a_{33})^{(6)}\left(\nu^{(6)}\right)^2 + (\sigma_1)^{(6)}\nu^{(6)} - (a_{32})^{(6)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(6)}, (\nu_0)^{(6)}$ :-

For $0 < \boxed{(\nu_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0}} < (\nu_1)^{(6)} < (\bar{\nu}_1)^{(6)}$

$$\nu^{(6)}(t) \geq \frac{(\nu_1)^{(6)}+(C)^{(6)}(\nu_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_0)^{(6)}\right)t\right]}}{1+(C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_0)^{(6)}\right)t\right]}} \quad , \quad \boxed{(C)^{(6)} = \frac{(\nu_1)^{(6)}-(\nu_0)^{(6)}}{(\nu_0)^{(6)}-(\nu_2)^{(6)}}}$$

it follows $(\nu_0)^{(6)} \leq \nu^{(6)}(t) \leq (\nu_1)^{(6)}$

In the same manner , we get

$$\nu^{(6)}(t) \leq \frac{(\bar{\nu}_1)^{(6)}+(\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}}{1+(\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(6)} = \frac{(\bar{\nu}_1)^{(6)}-(\nu_0)^{(6)}}{(\nu_0)^{(6)}-(\bar{\nu}_2)^{(6)}}}$$

From which we deduce $(\nu_0)^{(6)} \leq \nu^{(6)}(t) \leq (\bar{\nu}_1)^{(6)}$

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If $0 < (\nu_1)^{(6)} < (\nu_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0} < (\bar{\nu}_1)^{(6)}$ we find like in the previous case,

$$(\nu_1)^{(6)} \leq \frac{(\nu_1)^{(6)}+(C)^{(6)}(\nu_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_2)^{(6)}\right)t\right]}}{1+(C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_2)^{(6)}\right)t\right]}} \leq \nu^{(6)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(6)}+(\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}}{1+(\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}} \leq (\bar{\nu}_1)^{(6)}$$

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If $0 < (\nu_1)^{(6)} \leq (\bar{\nu}_1)^{(6)} \leq \boxed{(\nu_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0}}$ , we obtain

$$(\nu_1)^{(6)} \leq \nu^{(6)}(t) \leq \frac{(\bar{\nu}_1)^{(6)}+(\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}}{1+(\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}} \leq (\nu_0)^{(6)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(6)}(t)$ :-

$$(m_2)^{(6)} \leq \nu^{(6)}(t) \leq (m_1)^{(6)}, \quad \boxed{\nu^{(6)}(t) = \frac{G_{32}(t)}{G_{33}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(6)}(t)$ :-

$$(\mu_2)^{(6)} \leq u^{(6)}(t) \leq (\mu_1)^{(6)}, \quad \boxed{u^{(6)}(t) = \frac{T_{32}(t)}{T_{33}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{32}^{''})^{(6)} = (a_{33}^{''})^{(6)}, then\ (\sigma_1)^{(6)} = (\sigma_2)^{(6)}$ and in this case $(\nu_1)^{(6)} = (\bar{\nu}_1)^{(6)}$ if in addition $(\nu_0)^{(6)} = (\nu_1)^{(6)}$ then $\nu^{(6)}(t) = (\nu_0)^{(6)}$ and as a consequence $G_{32}(t) = (\nu_0)^{(6)}G_{33}(t)$ **this also defines** $(\nu_0)^{(6)}$ **for the special case** .

Analogously if $(b_{32}^{''})^{(6)} = (b_{33}^{''})^{(6)}, then\ (\tau_1)^{(6)} = (\tau_2)^{(6)}$ and then $(u_1)^{(6)} = (\bar{u}_1)^{(6)}$ if in addition $(u_0)^{(6)} = (u_1)^{(6)}$ then $T_{32}(t) = (u_0)^{(6)}T_{33}(t)$ This is an important consequence of the relation between $(\nu_1)^{(6)}$ and $(\bar{\nu}_1)^{(6)}$, **and definition of** $(u_0)^{(6)}$.

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**Proof** : From global equations we obtain

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$$\frac{dv^{(7)}}{dt} = (a_{36})^{(7)} - \left((a'_{36})^{(7)} - (a'_{37})^{(7)} + (a''_{36})^{(7)}(T_{37},t)\right) - (a''_{37})^{(7)}(T_{37},t)v^{(7)} - (a_{37})^{(7)}v^{(7)}$$

**Definition of** $v^{(7)}$ :- $\boxed{v^{(7)} = \frac{G_{36}}{G_{37}}}$

It follows

$$-\left((a_{37})^{(7)}\left(v^{(7)}\right)^2 + (\sigma_2)^{(7)}v^{(7)} - (a_{36})^{(7)}\right) \le \frac{dv^{(7)}}{dt} \le -\left((a_{37})^{(7)}\left(v^{(7)}\right)^2 + (\sigma_1)^{(7)}v^{(7)} - (a_{36})^{(7)}\right)$$

From which one obtains

**Definition of** $(\bar{v}_1)^{(7)}, (v_0)^{(7)}$ :-

For $0 < \boxed{(v_0)^{(7)} = \frac{G^0_{36}}{G^0_{37}}} < (v_1)^{(7)} < (\bar{v}_1)^{(7)}$

$$v^{(7)}(t) \ge \frac{(v_1)^{(7)} + (C)^{(7)}(v_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((v_1)^{(7)} - (v_0)^{(7)}\right)t\right]}}{1 + (C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((v_1)^{(7)} - (v_0)^{(7)}\right)t\right]}} \quad , \quad \boxed{(C)^{(7)} = \frac{(v_1)^{(7)} - (v_0)^{(7)}}{(v_0)^{(7)} - (v_2)^{(7)}}}$$

it follows $(v_0)^{(7)} \le v^{(7)}(t) \le (v_1)^{(7)}$

In the same manner , we get

$$v^{(7)}(t) \le \frac{(\bar{v}_1)^{(7)} + (\bar{C})^{(7)}(\bar{v}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}}{1 + (\bar{C})^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(7)} = \frac{(\bar{v}_1)^{(7)} - (v_0)^{(7)}}{(v_0)^{(7)} - (\bar{v}_2)^{(7)}}}$$

From which we deduce $(v_0)^{(7)} \le v^{(7)}(t) \le (\bar{v}_1)^{(7)}$

If $0 < (v_1)^{(7)} < (v_0)^{(7)} = \frac{G^0_{36}}{G^0_{37}} < (\bar{v}_1)^{(7)}$ we find like in the previous case,

$$(v_1)^{(7)} \le \frac{(v_1)^{(7)} + (C)^{(7)}(v_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((v_1)^{(7)} - (v_2)^{(7)}\right)t\right]}}{1 + (C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((v_1)^{(7)} - (v_2)^{(7)}\right)t\right]}} \le v^{(7)}(t) \le$$

$$\frac{(\bar{v}_1)^{(7)} + (\bar{C})^{(7)}(\bar{v}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}}{1 + (C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}} \le (\bar{v}_1)^{(7)}$$

If $0 < (v_1)^{(7)} \le (\bar{v}_1)^{(7)} \le \boxed{(v_0)^{(7)} = \frac{G^0_{36}}{G^0_{37}}}$ , we obtain

$$(v_1)^{(7)} \le v^{(7)}(t) \le \frac{(\bar{v}_1)^{(7)} + (\bar{C})^{(7)}(\bar{v}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}}{1 + (C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}} \le (v_0)^{(7)}$$

And so with the notation of the first part of condition (c) , we have

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**Definition of** $v^{(7)}(t)$ :-

$(m_2)^{(7)} \leq v^{(7)}(t) \leq (m_1)^{(7)}$, $\boxed{v^{(7)}(t) = \frac{G_{36}(t)}{G_{37}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(7)}(t)$ :-

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$(\mu_2)^{(7)} \leq u^{(7)}(t) \leq (\mu_1)^{(7)}$, $\boxed{u^{(7)}(t) = \frac{T_{36}(t)}{T_{37}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{36}^{''})^{(7)} = (a_{37}^{''})^{(7)}, then\ (\sigma_1)^{(7)} = (\sigma_2)^{(7)}$ and in this case $(\nu_1)^{(7)} = (\bar{\nu}_1)^{(7)}$ if in addition $(\nu_0)^{(7)} = (\nu_1)^{(7)}$ then $v^{(7)}(t) = (\nu_0)^{(7)}$ and as a consequence $G_{36}(t) = (\nu_0)^{(7)} G_{37}(t)$ **this also defines** $(\nu_0)^{(7)}$ **for the special case** .

Analogously if $(b_{36}^{''})^{(7)} = (b_{37}^{''})^{(7)}, then\ (\tau_1)^{(7)} = (\tau_2)^{(7)}$ and then $(u_1)^{(7)} = (\bar{u}_1)^{(7)}$if in addition $(u_0)^{(7)} = (u_1)^{(7)}$ then $T_{36}(t) = (u_0)^{(7)} T_{37}(t)$ This is an important consequence of the relation between $(\nu_1)^{(7)}$ and $(\bar{\nu}_1)^{(7)}$, **and definition of** $(u_0)^{(7)}$.

**Proof** : From global equations we obtain

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$$\frac{dv^{(8)}}{dt} = (a_{40})^{(8)} - \left((a_{40}^{'})^{(8)} - (a_{41}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41}, t)\right) - (a_{41}^{''})^{(8)}(T_{41}, t) v^{(8)} - (a_{41})^{(8)} v^{(8)}$$

**Definition of** $v^{(8)}$ :- $\boxed{v^{(8)} = \frac{G_{40}}{G_{41}}}$

It follows

$$-\left((a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_2)^{(8)} v^{(8)} - (a_{40})^{(8)}\right) \leq \frac{dv^{(8)}}{dt} \leq -\left((a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_1)^{(8)} v^{(8)} - (a_{40})^{(8)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(8)}, (\nu_0)^{(8)}$ :-

For $0 < \boxed{(\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}} < (\nu_1)^{(8)} < (\bar{\nu}_1)^{(8)}$

$$v^{(8)}(t) \geq \frac{(\nu_1)^{(8)} + (C)^{(8)}(\nu_2)^{(8)} e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}}{1 + (C)^{(8)} e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}} \quad , \quad \boxed{(C)^{(8)} = \frac{(\nu_1)^{(8)} - (\nu_0)^{(8)}}{(\nu_0)^{(8)} - (\nu_2)^{(8)}}}$$

it follows $(\nu_0)^{(8)} \leq \nu^{(8)}(t) \leq (\nu_1)^{(8)}$

In the same manner , we get 422

$$\nu^{(8)}(t) \leq \frac{(\bar{\nu}_1)^{(8)}+(\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}}{1+(\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(8)} = \frac{(\bar{\nu}_1)^{(8)}-(\nu_0)^{(8)}}{(\nu_0)^{(8)}-(\bar{\nu}_2)^{(8)}}}$$

From which we deduce $(\nu_0)^{(8)} \leq \nu^{(8)}(t) \leq (\bar{\nu}_8)^{(8)}$

If $0 < (\nu_1)^{(8)} < (\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0} < (\bar{\nu}_1)^{(8)}$ we find like in the previous case, 423

$$(\nu_1)^{(8)} \leq \frac{(\nu_1)^{(8)}+(C)^{(8)}(\nu_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)}-(\nu_2)^{(8)}\right)t\right]}}{1+(C)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)}-(\nu_2)^{(8)}\right)t\right]}} \leq \nu^{(8)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(8)}+(\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}}{1+(\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}} \leq (\bar{\nu}_1)^{(8)}$$

If $0 < (\nu_1)^{(8)} \leq (\bar{\nu}_1)^{(8)} \leq \boxed{(\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}}$ , we obtain 424

$$(\nu_1)^{(8)} \leq \nu^{(8)}(t) \leq \frac{(\bar{\nu}_1)^{(8)}+(C)^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}}{1+(C)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}} \leq (\nu_0)^{(8)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(8)}(t)$ :-

$$(m_2)^{(8)} \leq \nu^{(8)}(t) \leq (m_1)^{(8)}, \quad \boxed{\nu^{(8)}(t) = \frac{G_{40}(t)}{G_{41}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(8)}(t)$ :-

$$(\mu_2)^{(8)} \leq u^{(8)}(t) \leq (\mu_1)^{(8)}, \quad \boxed{u^{(8)}(t) = \frac{T_{40}(t)}{T_{41}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{40}^{''})^{(8)} = (a_{41}^{''})^{(8)}, then\ (\sigma_1)^{(8)} = (\sigma_2)^{(8)}$ and in this case $(\nu_1)^{(8)} = (\bar{\nu}_1)^{(8)}$ if in addition $(\nu_0)^{(8)} = (\nu_1)^{(8)}$ then $\nu^{(8)}(t) = (\nu_0)^{(8)}$ and as a consequence $G_{40}(t) = (\nu_0)^{(8)}G_{41}(t)$ **this also defines** $(\nu_0)^{(8)}$ **for the special case** .

Analogously if $(b_{40}^{''})^{(8)} = (b_{41}^{''})^{(8)}, then\ (\tau_1)^{(8)} = (\tau_2)^{(8)}$ and then $(u_1)^{(8)} = (\bar{u}_1)^{(8)}$ if in addition $(u_0)^{(8)} = (u_1)^{(8)}$ then $T_{40}(t) = (u_0)^{(8)}T_{41}(t)$ This is an important

consequence of the relation between $(\nu_1)^{(8)}$ and $(\bar{\nu}_1)^{(8)}$, **and definition of** $(u_0)^{(8)}$.

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**Proof** : From 99,20,44,22,23,44 we obtain

$$\frac{d\nu^{(9)}}{dt} = (a_{44})^{(9)} - \left((a'_{44})^{(9)} - (a'_{45})^{(9)} + (a''_{44})^{(9)}(T_{45},t)\right) - (a''_{45})^{(9)}(T_{45},t)\nu^{(9)} - (a_{45})^{(9)}\nu^{(9)}$$

**Definition of** $\nu^{(9)}$ :- $\boxed{\nu^{(9)} = \frac{G_{44}}{G_{45}}}$

It follows

$$-\left((a_{45})^{(9)}\left(\nu^{(9)}\right)^2 + (\sigma_2)^{(9)}\nu^{(9)} - (a_{44})^{(9)}\right) \le \frac{d\nu^{(9)}}{dt} \le -\left((a_{45})^{(9)}\left(\nu^{(9)}\right)^2 + (\sigma_1)^{(9)}\nu^{(9)} - (a_{44})^{(9)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(9)}, (\nu_0)^{(9)}$ :-

For $0 < \boxed{(\nu_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}}} < (\nu_1)^{(9)} < (\bar{\nu}_1)^{(9)}$

$$\nu^{(9)}(t) \ge \frac{(\nu_1)^{(9)} + (C)^{(9)}(\nu_2)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_0)^{(9)}\right)t\right]}}{1 + (C)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_0)^{(9)}\right)t\right]}} \quad , \quad \boxed{(C)^{(9)} = \frac{(\nu_1)^{(9)} - (\nu_0)^{(9)}}{(\nu_0)^{(9)} - (\nu_2)^{(9)}}}$$

it follows $(\nu_0)^{(9)} \le \nu^{(9)}(t) \le (\nu_9)^{(9)}$

In the same manner , we get

$$\nu^{(9)}(t) \le \frac{(\bar{\nu}_1)^{(9)} + (\bar{C})^{(9)}(\bar{\nu}_2)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}}{1 + (\bar{C})^{(9)} e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(9)} = \frac{(\bar{\nu}_1)^{(9)} - (\nu_0)^{(9)}}{(\nu_0)^{(9)} - (\bar{\nu}_2)^{(9)}}}$$

From which we deduce $(\nu_0)^{(9)} \le \nu^{(9)}(t) \le (\bar{\nu}_1)^{(9)}$

If $0 < (\nu_1)^{(9)} < (\nu_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}} < (\bar{\nu}_1)^{(9)}$ we find like in the previous case,

$$(\nu_1)^{(9)} \le \frac{(\nu_1)^{(9)} + (C)^{(9)}(\nu_2)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_2)^{(9)}\right)t\right]}}{1 + (C)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_2)^{(9)}\right)t\right]}} \le \nu^{(9)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(9)} + (\bar{C})^{(9)}(\bar{\nu}_2)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}}{1 + (\bar{C})^{(9)} e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}} \le (\bar{\nu}_1)^{(9)}$$

If $0 < (\nu_1)^{(9)} \le (\bar{\nu}_1)^{(9)} \le \boxed{(\nu_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}}}$ , we obtain

$$(\nu_1)^{(9)} \le \nu^{(9)}(t) \le \frac{(\bar{\nu}_1)^{(9)} + (\bar{C})^{(9)}(\bar{\nu}_2)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}}{1 + (\bar{C})^{(9)} e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}} \le (\nu_0)^{(9)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(9)}(t)$ :-

$(m_2)^{(9)} \leq \nu^{(9)}(t) \leq (m_1)^{(9)}$, $\boxed{\nu^{(9)}(t) = \frac{G_{44}(t)}{G_{45}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(9)}(t)$ :-

$(\mu_2)^{(9)} \leq u^{(9)}(t) \leq (\mu_1)^{(9)}$, $\boxed{u^{(9)}(t) = \frac{T_{44}(t)}{T_{45}(t)}}$

Now, using this result and replacing it in 99, 20,44,22,23, and 44 we get easily the result stated in the theorem.

**Particular case :**

If $(a_{44}^{''})^{(9)} = (a_{45}^{''})^{(9)}, then\ (\sigma_1)^{(9)} = (\sigma_2)^{(9)}$ and in this case $(\nu_1)^{(9)} = (\bar{\nu}_1)^{(9)}$ if in addition $(\nu_0)^{(9)} = (\nu_1)^{(9)}$ then $\nu^{(9)}(t) = (\nu_0)^{(9)}$ and as a consequence $G_{44}(t) = (\nu_0)^{(9)} G_{45}(t)$ **this also defines** $(\nu_0)^{(9)}$ **for the special case** .

Analogously if $(b_{44}^{''})^{(9)} = (b_{45}^{''})^{(9)}, then\ (\tau_1)^{(9)} = (\tau_2)^{(9)}$ and then $(u_1)^{(9)} = (\bar{u}_1)^{(9)}$if in addition $(u_0)^{(9)} = (u_1)^{(9)}$ then $T_{44}(t) = (u_0)^{(9)} T_{45}(t)$ This is an important consequence of the relation between $(\nu_1)^{(9)}$ and $(\bar{\nu}_1)^{(9)}$, **and definition of** $(u_0)^{(9)}$.

We can prove the following 425

**Theorem :** If $(a_i^{''})^{(1)} and\ (b_i^{''})^{(1)}$ are independent on $t$ , and the conditions with the notations

$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} < 0$

$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a_{13})^{(1)}(p_{13})^{(1)} + (a_{14}^{'})^{(1)}(p_{14})^{(1)} + (p_{13})^{(1)}(p_{14})^{(1)} > 0$

$(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} > 0$ ,

$(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} - (b_{13}^{'})^{(1)}(r_{14})^{(1)} - (b_{14}^{'})^{(1)}(r_{14})^{(1)} + (r_{13})^{(1)}(r_{14})^{(1)} < 0$

$with\ (p_{13})^{(1)}, (r_{14})^{(1)}$ as defined by equation are satisfied , then the system

**Theorem :** If $(a_i^{''})^{(2)} and\ (b_i^{''})^{(2)}$ are independent on t , and the conditions with the notations 426

$(a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} < 0$ 427

$(a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a_{16})^{(2)}(p_{16})^{(2)} + (a_{17}^{'})^{(2)}(p_{17})^{(2)} + (p_{16})^{(2)}(p_{17})^{(2)} > 0$ 428

$(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} > 0$ , 429

$(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} - (b_{16}^{'})^{(2)}(r_{17})^{(2)} - (b_{17}^{'})^{(2)}(r_{17})^{(2)} + (r_{16})^{(2)}(r_{17})^{(2)} < 0$ 430

$with\ (p_{16})^{(2)}, (r_{17})^{(2)}$ as defined by equation are satisfied , then the system

**Theorem :** If $(a_i^{''})^{(3)} and\ (b_i^{''})^{(3)}$ are independent on $t$ , and the conditions with the notations 431

$(a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} < 0$

| | |
|---|---|
| $(a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a_{20})^{(3)}(p_{20})^{(3)} + (a_{21}^{'})^{(3)}(p_{21})^{(3)} + (p_{20})^{(3)}(p_{21})^{(3)} > 0$<br>$(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} > 0$ ,<br>$(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} - (b_{20}^{'})^{(3)}(r_{21})^{(3)} - (b_{21}^{'})^{(3)}(r_{21})^{(3)} + (r_{20})^{(3)}(r_{21})^{(3)} < 0$<br>$with\ \ (p_{20})^{(3)}, (r_{21})^{(3)}$ as defined by equation are satisfied , then the system | |
| We can prove the following<br>**Theorem :** If $(a_i^{''})^{(4)} and\ (b_i^{''})^{(4)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} < 0$<br>$(a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a_{24})^{(4)}(p_{24})^{(4)} + (a_{25}^{'})^{(4)}(p_{25})^{(4)} + (p_{24})^{(4)}(p_{25})^{(4)} > 0$<br>$(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} > 0$ ,<br>$(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} - (b_{24}^{'})^{(4)}(r_{25})^{(4)} - (b_{25}^{'})^{(4)}(r_{25})^{(4)} + (r_{24})^{(4)}(r_{25})^{(4)} < 0$<br>$with\ \ (p_{24})^{(4)}, (r_{25})^{(4)}$ as defined by equation are satisfied , then the system | 432 |
| **Theorem :** If $(a_i^{''})^{(5)} and\ (b_i^{''})^{(5)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} < 0$<br>$(a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a_{28})^{(5)}(p_{28})^{(5)} + (a_{29}^{'})^{(5)}(p_{29})^{(5)} + (p_{28})^{(5)}(p_{29})^{(5)} > 0$<br>$(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} > 0$ ,<br>$(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} - (b_{28}^{'})^{(5)}(r_{29})^{(5)} - (b_{29}^{'})^{(5)}(r_{29})^{(5)} + (r_{28})^{(5)}(r_{29})^{(5)} < 0$<br>$with\ \ (p_{28})^{(5)}, (r_{29})^{(5)}$ as defined by equation are satisfied , then the system | 433 |
| **Theorem** If $(a_i^{''})^{(6)} and\ (b_i^{''})^{(6)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} < 0$<br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32})^{(6)}(p_{32})^{(6)} + (a_{33}^{'})^{(6)}(p_{33})^{(6)} + (p_{32})^{(6)}(p_{33})^{(6)} > 0$<br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} > 0$ ,<br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} - (b_{32}^{'})^{(6)}(r_{33})^{(6)} - (b_{33}^{'})^{(6)}(r_{33})^{(6)} + (r_{32})^{(6)}(r_{33})^{(6)} < 0$<br>$with\ \ (p_{32})^{(6)}, (r_{33})^{(6)}$ as defined by equation are satisfied , then the system | 434 |
| **Theorem :** If $(a_i^{''})^{(7)} and\ (b_i^{''})^{(7)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} < 0$ | 435 |

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| $(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36})^{(7)}(p_{36})^{(7)} + (a_{37}^{'})^{(7)}(p_{37})^{(7)} + (p_{36})^{(7)}(p_{37})^{(7)} > 0$<br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} > 0$ ,<br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} - (b_{36}^{'})^{(7)}(r_{37})^{(7)} - (b_{37}^{'})^{(7)}(r_{37})^{(7)} + (r_{36})^{(7)}(r_{37})^{(7)} < 0$<br>$with\ \ (p_{36})^{(7)}, (r_{37})^{(7)}$ as defined by equation are satisfied , then the system | |
| **Theorem :** If $(a_i^{''})^{(8)} and\ (b_i^{''})^{(8)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} < 0$<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40})^{(8)}(p_{40})^{(8)} + (a_{41}^{'})^{(8)}(p_{41})^{(8)} + (p_{40})^{(8)}(p_{41})^{(8)} > 0$<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} > 0$ ,<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} - (b_{40}^{'})^{(8)}(r_{41})^{(8)} - (b_{41}^{'})^{(8)}(r_{41})^{(8)} + (r_{40})^{(8)}(r_{41})^{(8)} < 0$<br>$with\ \ (p_{40})^{(8)}, (r_{41})^{(8)}$ as defined by equation are satisfied , then the system | 436 |
| **Theorem :** If $(a_i^{''})^{(9)} and\ (b_i^{''})^{(9)}$ are independent on $t$ , and the conditions (with the notations 45,46,27,28)<br>$(a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} < 0$<br>$(a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44})^{(9)}(p_{44})^{(9)} + (a_{45}^{'})^{(9)}(p_{45})^{(9)} + (p_{44})^{(9)}(p_{45})^{(9)} > 0$<br>$(b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} > 0$ ,<br>$(b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} - (b_{44}^{'})^{(9)}(r_{45})^{(9)} - (b_{45}^{'})^{(9)}(r_{45})^{(9)} + (r_{44})^{(9)}(r_{45})^{(9)} < 0$<br>$with\ \ (p_{44})^{(9)}, (r_{45})^{(9)}$ as defined by equation 45 are satisfied , then the system | 436 A |
| $(a_{13})^{(1)}G_{14} - \left[(a_{13}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14})\right]G_{13} = 0$ | 437 |
| $(a_{14})^{(1)}G_{13} - \left[(a_{14}^{'})^{(1)} + (a_{14}^{''})^{(1)}(T_{14})\right]G_{14} = 0$ | 438 |
| $(a_{15})^{(1)}G_{14} - \left[(a_{15}^{'})^{(1)} + (a_{15}^{''})^{(1)}(T_{14})\right]G_{15} = 0$ | 439 |
| $(b_{13})^{(1)}T_{14} - [(b_{13}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G)\ ]T_{13} = 0$ | 440 |
| $(b_{14})^{(1)}T_{13} - [(b_{14}^{'})^{(1)} - (b_{14}^{''})^{(1)}(G)\ ]T_{14} = 0$ | 441 |
| $(b_{15})^{(1)}T_{14} - [(b_{15}^{'})^{(1)} - (b_{15}^{''})^{(1)}(G)\ ]T_{15} = 0$ | 442 |
| has a unique positive solution , which is an equilibrium solution for the system | |

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| $(a_{16})^{(2)}G_{17} - \left[(a_{16}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17})\right]G_{16} = 0$ | 443 |
| $(a_{17})^{(2)}G_{16} - \left[(a_{17}^{'})^{(2)} + (a_{17}^{''})^{(2)}(T_{17})\right]G_{17} = 0$ | 444 |
| $(a_{18})^{(2)}G_{17} - \left[(a_{18}^{'})^{(2)} + (a_{18}^{''})^{(2)}(T_{17})\right]G_{18} = 0$ | 445 |
| $(b_{16})^{(2)}T_{17} - [(b_{16}^{'})^{(2)} - (b_{16}^{''})^{(2)}(G_{19})\,]T_{16} = 0$ | 446 |
| $(b_{17})^{(2)}T_{16} - [(b_{17}^{'})^{(2)} - (b_{17}^{''})^{(2)}(G_{19})\,]T_{17} = 0$ | 447 |
| $(b_{18})^{(2)}T_{17} - [(b_{18}^{'})^{(2)} - (b_{18}^{''})^{(2)}(G_{19})\,]T_{18} = 0$ | 448 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{20})^{(3)}G_{21} - \left[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21})\right]G_{20} = 0$ | 449 |
| $(a_{21})^{(3)}G_{20} - \left[(a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)}(T_{21})\right]G_{21} = 0$ | 450 |
| $(a_{22})^{(3)}G_{21} - \left[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21})\right]G_{22} = 0$ | 451 |
| $(b_{20})^{(3)}T_{21} - [(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23})\,]T_{20} = 0$ | 452 |
| $(b_{21})^{(3)}T_{20} - [(b_{21}^{'})^{(3)} - (b_{21}^{''})^{(3)}(G_{23})\,]T_{21} = 0$ | 453 |
| $(b_{22})^{(3)}T_{21} - [(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23})\,]T_{22} = 0$ | 454 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25})\right]G_{24} = 0$ | 455 |
| $(a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25})\right]G_{25} = 0$ | 456 |
| $(a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25})\right]G_{26} = 0$ | 457 |
| $(b_{24})^{(4)}T_{25} - [(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}\left((G_{27})\right)\,]T_{24} = 0$ | 458 |
| $(b_{25})^{(4)}T_{24} - [(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}\left((G_{27})\right)\,]T_{25} = 0$ | 459 |
| $(b_{26})^{(4)}T_{25} - [(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}\left((G_{27})\right)\,]T_{26} = 0$ | 460 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29})\right]G_{28} = 0$ | 461 |
| $(a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29})\right]G_{29} = 0$ | 462 |

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| $(a_{30})^{(5)}G_{29} - \left[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29})\right]G_{30} = 0$ | 463 |
| $(b_{28})^{(5)}T_{29} - [(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}(G_{31})\,]T_{28} = 0$ | 464 |
| $(b_{29})^{(5)}T_{28} - [(b_{29}^{'})^{(5)} - (b_{29}^{''})^{(5)}(G_{31})\,]T_{29} = 0$ | 465 |
| $(b_{30})^{(5)}T_{29} - [(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}(G_{31})\,]T_{30} = 0$ | 466 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{32})^{(6)}G_{33} - \left[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33})\right]G_{32} = 0$ | 467 |
| $(a_{33})^{(6)}G_{32} - \left[(a_{33}^{'})^{(6)} + (a_{33}^{''})^{(6)}(T_{33})\right]G_{33} = 0$ | 468 |
| $(a_{34})^{(6)}G_{33} - \left[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33})\right]G_{34} = 0$ | 469 |
| $(b_{32})^{(6)}T_{33} - [(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}(G_{35})\,]T_{32} = 0$ | 470 |
| $(b_{33})^{(6)}T_{32} - [(b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}(G_{35})\,]T_{33} = 0$ | 471 |
| $(b_{34})^{(6)}T_{33} - [(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}(G_{35})\,]T_{34} = 0$ | 472 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{36})^{(7)}G_{37} - \left[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37})\right]G_{36} = 0$ | 473 |
| $(a_{37})^{(7)}G_{36} - \left[(a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37})\right]G_{37} = 0$ | 474 |
| $(a_{38})^{(7)}G_{37} - \left[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37})\right]G_{38} = 0$ | 475 |
| $(b_{36})^{(7)}T_{37} - [(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}(G_{39})]T_{36} = 0$ | 476 |
| $(b_{37})^{(7)}T_{36} - [(b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}(G_{39})\,]T_{37} = 0$ | 477 |
| $(b_{38})^{(7)}T_{37} - [(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}(G_{39})\,]T_{38} = 0$ | 478 |
| $(a_{40})^{(8)}G_{41} - \left[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41})\right]G_{40} = 0$ | 479 |
| $(a_{41})^{(8)}G_{40} - \left[(a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41})\right]G_{41} = 0$ | 480 |

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| $(a_{42})^{(8)}G_{41} - \left[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41})\right]G_{42} = 0$ | 481 |
| $(b_{40})^{(8)}T_{41} - [(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}(G_{43})\,]T_{40} = 0$ | 482 |
| $(b_{41})^{(8)}T_{40} - [(b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}(G_{43})\,]T_{41} = 0$ | 483 |
| $(b_{42})^{(8)}T_{41} - [(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}(G_{43})\,]T_{42} = 0$ | 484 |
| $(a_{44})^{(9)}G_{45} - \left[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45})\right]G_{44} = 0$ | 484 A |
| $(a_{45})^{(9)}G_{44} - \left[(a_{45}^{'})^{(9)} + (a_{45}^{''})^{(9)}(T_{45})\right]G_{45} = 0$ | |
| $(a_{46})^{(9)}G_{45} - \left[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45})\right]G_{46} = 0$ | |
| $(b_{44})^{(9)}T_{45} - [(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}(G_{47})\,]T_{44} = 0$ | |
| $(b_{45})^{(9)}T_{44} - [(b_{45}^{'})^{(9)} - (b_{45}^{''})^{(9)}(G_{47})\,]T_{45} = 0$ | |
| $(b_{46})^{(9)}T_{45} - [(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}(G_{47})\,]T_{46} = 0$ | |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{13}, G_{14}$ if<br><br>$F(T) = (a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a_{13}^{'})^{(1)}(a_{14}^{''})^{(1)}(T_{14}) + (a_{14}^{'})^{(1)}(a_{13}^{''})^{(1)}(T_{14}) + (a_{13}^{''})^{(1)}(T_{14})(a_{14}^{''})^{(1)}(T_{14}) = 0$ | 485 |
| **Proof:**<br><br>(f) Indeed the first two equations have a nontrivial solution $G_{16}, G_{17}$ if<br><br>$\mathrm{F}(T_{19}) = (a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a_{16}^{'})^{(2)}(a_{17}^{''})^{(2)}(T_{17}) + (a_{17}^{'})^{(2)}(a_{16}^{''})^{(2)}(T_{17}) + (a_{16}^{''})^{(2)}(T_{17})(a_{17}^{''})^{(2)}(T_{17}) = 0$ | 486 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{20}, G_{21}$ if<br><br>$F(T_{23}) = (a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a_{20}^{'})^{(3)}(a_{21}^{''})^{(3)}(T_{21}) + (a_{21}^{'})^{(3)}(a_{20}^{''})^{(3)}(T_{21}) + (a_{20}^{''})^{(3)}(T_{21})(a_{21}^{''})^{(3)}(T_{21}) = 0$ | 487 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{24}, G_{25}$ if<br><br>$F(T_{27}) = (a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a_{24}^{'})^{(4)}(a_{25}^{''})^{(4)}(T_{25}) + (a_{25}^{'})^{(4)}(a_{24}^{''})^{(4)}(T_{25}) +$ | 488 |

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| $(a_{24}^{''})^{(4)}(T_{25})(a_{25}^{''})^{(4)}(T_{25}) = 0$ | |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{28}, G_{29}$ if<br><br>$F(T_{31}) = (a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a_{28}^{'})^{(5)}(a_{29}^{''})^{(5)}(T_{29}) + (a_{29}^{'})^{(5)}(a_{28}^{''})^{(5)}(T_{29}) + (a_{28}^{''})^{(5)}(T_{29})(a_{29}^{''})^{(5)}(T_{29}) = 0$ | 489 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{32}, G_{33}$ if<br><br>$F(T_{35}) = (a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32}^{'})^{(6)}(a_{33}^{''})^{(6)}(T_{33}) + (a_{33}^{'})^{(6)}(a_{32}^{''})^{(6)}(T_{33}) + (a_{32}^{''})^{(6)}(T_{33})(a_{33}^{''})^{(6)}(T_{33}) = 0$ | 490 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{36}, G_{37}$ if<br><br>$F(T_{39}) = (a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36}^{'})^{(7)}(a_{37}^{''})^{(7)}(T_{37}) + (a_{37}^{'})^{(7)}(a_{36}^{''})^{(7)}(T_{37}) + (a_{36}^{''})^{(7)}(T_{37})(a_{37}^{''})^{(7)}(T_{37}) = 0$ | 491 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{40}, G_{41}$ if<br><br>$F(T_{43}) = (a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40}^{'})^{(8)}(a_{41}^{''})^{(8)}(T_{41}) + (a_{41}^{'})^{(8)}(a_{40}^{''})^{(8)}(T_{41}) + (a_{40}^{''})^{(8)}(T_{41})(a_{41}^{''})^{(8)}(T_{41}) = 0$ | 492 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{44}, G_{45}$ if<br><br>$F(T_{47}) = (a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44}^{'})^{(9)}(a_{45}^{''})^{(9)}(T_{45}) + (a_{45}^{'})^{(9)}(a_{44}^{''})^{(9)}(T_{45}) + (a_{44}^{''})^{(9)}(T_{45})(a_{45}^{''})^{(9)}(T_{45}) = 0$ | 492 A |
| **Definition and uniqueness of** $T_{14}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{14})$ being increasing, it follows that there exists a unique $T_{14}^{*}$ for which $f(T_{14}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{13} = \frac{(a_{13})^{(1)}G_{14}}{[(a_{13}^{'})^{(1)}+(a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15} = \frac{(a_{15})^{(1)}G_{14}}{[(a_{15}^{'})^{(1)}+(a_{15}^{''})^{(1)}(T_{14}^{*})]}$ | 493 |
| **Definition and uniqueness of** $T_{17}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(2)}(T_{17})$ being increasing, it follows that there exists a unique $T_{17}^{*}$ for which $f(T_{17}^{*}) = 0$. With this value , we obtain from the three first | 494 |

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| equations | |
| $G_{16} = \frac{(a_{16})^{(2)}G_{17}}{[(a_{16}^{'})^{(2)}+(a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18} = \frac{(a_{18})^{(2)}G_{17}}{[(a_{18}^{'})^{(2)}+(a_{18}^{''})^{(2)}(T_{17}^{*})]}$ | 495 |
| **Definition and uniqueness of** $T_{21}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{21})$ being increasing, it follows that there exists a unique $T_{21}^{*}$ for which $f(T_{21}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{20} = \frac{(a_{20})^{(3)}G_{21}}{[(a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21}^{*})]}$ , $G_{22} = \frac{(a_{22})^{(3)}G_{21}}{[(a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21}^{*})]}$ | 496 |
| **Definition and uniqueness of** $T_{25}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(4)}(T_{25})$ being increasing, it follows that there exists a unique $T_{25}^{*}$ for which $f(T_{25}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{24} = \frac{(a_{24})^{(4)}G_{25}}{[(a_{24}^{'})^{(4)}+(a_{24}^{''})^{(4)}(T_{25}^{*})]}$ , $G_{26} = \frac{(a_{26})^{(4)}G_{25}}{[(a_{26}^{'})^{(4)}+(a_{26}^{''})^{(4)}(T_{25}^{*})]}$ | 497 |
| **Definition and uniqueness of** $T_{29}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(5)}(T_{29})$ being increasing, it follows that there exists a unique $T_{29}^{*}$ for which $f(T_{29}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{28} = \frac{(a_{28})^{(5)}G_{29}}{[(a_{28}^{'})^{(5)}+(a_{28}^{''})^{(5)}(T_{29}^{*})]}$ , $G_{30} = \frac{(a_{30})^{(5)}G_{29}}{[(a_{30}^{'})^{(5)}+(a_{30}^{''})^{(5)}(T_{29}^{*})]}$ | 498 |
| **Definition and uniqueness of** $T_{33}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(6)}(T_{33})$ being increasing, it follows that there exists a unique $T_{33}^{*}$ for which $f(T_{33}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{32} = \frac{(a_{32})^{(6)}G_{33}}{[(a_{32}^{'})^{(6)}+(a_{32}^{''})^{(6)}(T_{33}^{*})]}$ , $G_{34} = \frac{(a_{34})^{(6)}G_{33}}{[(a_{34}^{'})^{(6)}+(a_{34}^{''})^{(6)}(T_{33}^{*})]}$ | 499 |
| **Definition and uniqueness of** $T_{37}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(7)}(T_{37})$ being increasing, it follows that there exists a unique $T_{37}^{*}$ for which $f(T_{37}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{36} = \frac{(a_{36})^{(7)}G_{37}}{[(a_{36}^{'})^{(7)}+(a_{36}^{''})^{(7)}(T_{37}^{*})]}$ , $G_{38} = \frac{(a_{38})^{(7)}G_{37}}{[(a_{38}^{'})^{(7)}+(a_{38}^{''})^{(7)}(T_{37}^{*})]}$ | 500 |

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| **Definition and uniqueness of** $T_{41}^{*}$ :-<br><br>After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(8)}(T_{41})$ being increasing, it follows that there exists a unique $T_{41}^{*}$ for which $f(T_{41}^{*})=0$. With this value , we obtain from the three first equations<br><br>$G_{40}=\frac{(a_{40})^{(8)}G_{41}}{[(a_{40}^{'})^{(8)}+(a_{40}^{''})^{(8)}(T_{41}^{*})]}$ , $G_{42}=\frac{(a_{42})^{(8)}G_{41}}{[(a_{42}^{'})^{(8)}+(a_{42}^{''})^{(8)}(T_{41}^{*})]}$ | 501 |
| **Definition and uniqueness of** $T_{45}^{*}$ :-<br><br>After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(9)}(T_{45})$ being increasing, it follows that there exists a unique $T_{45}^{*}$ for which $f(T_{45}^{*})=0$. With this value , we obtain from the three first equations<br><br>$G_{44}=\frac{(a_{44})^{(9)}G_{45}}{[(a_{44}^{'})^{(9)}+(a_{44}^{''})^{(9)}(T_{45}^{*})]}$ , $G_{46}=\frac{(a_{46})^{(9)}G_{45}}{[(a_{46}^{'})^{(9)}+(a_{46}^{''})^{(9)}(T_{45}^{*})]}$ | 501 A |
| By the same argument, the equations admit solutions $G_{13}, G_{14}$ if<br>$\varphi(G)=(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)}-(b_{13})^{(1)}(b_{14})^{(1)}-$<br>$[(b_{13}^{'})^{(1)}(b_{14}^{''})^{(1)}(G)+(b_{14}^{'})^{(1)}(b_{13}^{''})^{(1)}(G)]+(b_{13}^{''})^{(1)}(G)(b_{14}^{''})^{(1)}(G)=0$<br><br>Where in $G(G_{13},G_{14},G_{15}), G_{13}, G_{15}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{14}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{14}^{*}$ such that $\varphi(G^{*})=0$ | 502 |
| By the same argument, the equations admit solutions $G_{16}, G_{17}$ if<br><br>$\varphi(G_{19})=(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)}-(b_{16})^{(2)}(b_{17})^{(2)}-$<br>$[(b_{16}^{'})^{(2)}(b_{17}^{''})^{(2)}(G_{19})+(b_{17}^{'})^{(2)}(b_{16}^{''})^{(2)}(G_{19})]+(b_{16}^{''})^{(2)}(G_{19})(b_{17}^{''})^{(2)}(G_{19})=0$ | 503 |
| Where in $(G_{19})(G_{16},G_{17},G_{18}), G_{16}, G_{18}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{17}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{14}^{*}$ such that $\varphi((G_{19})^{*})=0$ | 504 |
| By the same argument, the equations admit solutions $G_{20}, G_{21}$ if<br><br>$\varphi(G_{23})=(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)}-(b_{20})^{(3)}(b_{21})^{(3)}-$<br>$[(b_{20}^{'})^{(3)}(b_{21}^{''})^{(3)}(G_{23})+(b_{21}^{'})^{(3)}(b_{20}^{''})^{(3)}(G_{23})]+(b_{20}^{''})^{(3)}(G_{23})(b_{21}^{''})^{(3)}(G_{23})=0$<br><br>Where in $G_{23}(G_{20},G_{21},G_{22}), G_{20}, G_{22}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{21}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{21}^{*}$ such that $\varphi((G_{23})^{*})=0$ | 505 |

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| By the same argument, the equations admit solutions $G_{24}, G_{25}$ if<br><br>$\varphi(G_{27}) = (b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} -$<br><br>$\left[(b_{24}^{'})^{(4)}(b_{25}^{''})^{(4)}(G_{27}) + (b_{25}^{'})^{(4)}(b_{24}^{''})^{(4)}(G_{27})\right] + (b_{24}^{''})^{(4)}(G_{27})(b_{25}^{''})^{(4)}(G_{27}) = 0$<br><br>Where in $(G_{27})(G_{24}, G_{25}, G_{26}), G_{24}, G_{26}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{25}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{25}^{*}$ such that $\varphi((G_{27})^{*}) = 0$ | 506 |
| By the same argument, the equations admit solutions $G_{28}, G_{29}$ if<br>$\varphi(G_{31}) = (b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} -$<br><br>$\left[(b_{28}^{'})^{(5)}(b_{29}^{''})^{(5)}(G_{31}) + (b_{29}^{'})^{(5)}(b_{28}^{''})^{(5)}(G_{31})\right] + (b_{28}^{''})^{(5)}(G_{31})(b_{29}^{''})^{(5)}(G_{31}) = 0$<br><br>Where in $(G_{31})(G_{28}, G_{29}, G_{30}), G_{28}, G_{30}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{29}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{29}^{*}$ such that $\varphi((G_{31})^{*}) = 0$ | 507 |
| By the same argument, the equations admit solutions $G_{32}, G_{33}$ if<br>$\varphi(G_{35}) = (b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} -$<br><br>$\left[(b_{32}^{'})^{(6)}(b_{33}^{''})^{(6)}(G_{35}) + (b_{33}^{'})^{(6)}(b_{32}^{''})^{(6)}(G_{35})\right] + (b_{32}^{''})^{(6)}(G_{35})(b_{33}^{''})^{(6)}(G_{35}) = 0$<br><br>Where in $(G_{35})(G_{32}, G_{33}, G_{34}), G_{32}, G_{34}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{33}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{33}^{*}$ such that $\varphi({G_{35}}^{*}) = 0$ | 508 |
| By the same argument, the equations admit solutions $G_{36}, G_{37}$ if<br>$\varphi(G_{39}) = (b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} -$<br>$\left[(b_{36}^{'})^{(7)}(b_{37}^{''})^{(7)}(G_{39}) + (b_{37}^{'})^{(7)}(b_{36}^{''})^{(7)}(G_{39})\right] + (b_{36}^{''})^{(7)}(G_{39})(b_{37}^{''})^{(7)}(G_{39}) = 0$<br><br>Where in $(G_{39})(G_{36}, G_{37}, G_{38}), G_{36}, G_{38}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{37}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{37}^{*}$ such that $\varphi({G_{39}}^{*}) = 0$ | 509 |
| By the same argument, the equations admit solutions $G_{40}, G_{41}$ if<br>$\varphi(G_{43}) = (b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} -$<br>$\left[(b_{40}^{'})^{(8)}(b_{41}^{''})^{(8)}(G_{43}) + (b_{41}^{'})^{(8)}(b_{40}^{''})^{(8)}(G_{43})\right] + (b_{40}^{''})^{(8)}(G_{43})(b_{41}^{''})^{(8)}(G_{43}) = 0$<br><br>Where in $(G_{43})(G_{40}, G_{41}, G_{42}), G_{40}, G_{42}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{41}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{41}^{*}$ such that $\varphi({G_{43}}^{*}) = 0$ | 510 |
| By the same argument, the equations 92,93 admit solutions $G_{44}, G_{45}$ if<br>$\varphi(G_{47}) = (b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} -$ | |

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| $\left[(b_{44}^{'})^{(9)}(b_{45}^{''})^{(9)}(G_{47})+(b_{45}^{'})^{(9)}(b_{44}^{''})^{(9)}(G_{47})\right]+(b_{44}^{''})^{(9)}(G_{47})(b_{45}^{''})^{(9)}(G_{47})=0$<br><br>Where in $(G_{47})(G_{44},G_{45},G_{46}),G_{44},G_{46}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{45}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{45}^{*}$ such that $\varphi((G_{47})^{*})=0$ | |
| Finally we obtain the unique solution<br><br>$G_{14}^{*}$ given by $\varphi(G^{*})=0$ , $T_{14}^{*}$ given by $f(T_{14}^{*})=0$ and<br><br>$G_{13}^{*}=\frac{(a_{13})^{(1)}G_{14}^{*}}{[(a_{13}^{'})^{(1)}+(a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15}^{*}=\frac{(a_{15})^{(1)}G_{14}^{*}}{[(a_{15}^{'})^{(1)}+(a_{15}^{''})^{(1)}(T_{14}^{*})]}$<br><br>$T_{13}^{*}=\frac{(b_{13})^{(1)}T_{14}^{*}}{[(b_{13}^{'})^{(1)}-(b_{13}^{''})^{(1)}(G^{*})]}$ , $T_{15}^{*}=\frac{(b_{15})^{(1)}T_{14}^{*}}{[(b_{15}^{'})^{(1)}-(b_{15}^{''})^{(1)}(G^{*})]}$<br><br>Obviously, these values represent an equilibrium solution | 511 |
| Finally we obtain the unique solution | |
| $G_{17}^{*}$ given by $\varphi((G_{19})^{*})=0$ , $T_{17}^{*}$ given by $f(T_{17}^{*})=0$ and | 512 |
| $G_{16}^{*}=\frac{(a_{16})^{(2)}G_{17}^{*}}{[(a_{16}^{'})^{(2)}+(a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18}^{*}=\frac{(a_{18})^{(2)}G_{17}^{*}}{[(a_{18}^{'})^{(2)}+(a_{18}^{''})^{(2)}(T_{17}^{*})]}$ | 513 |
| $T_{16}^{*}=\frac{(b_{16})^{(2)}T_{17}^{*}}{[(b_{16}^{'})^{(2)}-(b_{16}^{''})^{(2)}((G_{19})^{*})]}$ , $T_{18}^{*}=\frac{(b_{18})^{(2)}T_{17}^{*}}{[(b_{18}^{'})^{(2)}-(b_{18}^{''})^{(2)}((G_{19})^{*})]}$ | 514 |
| Obviously, these values represent an equilibrium solution | |
| Finally we obtain the unique solution<br><br>$G_{21}^{*}$ given by $\varphi((G_{23})^{*})=0$ , $T_{21}^{*}$ given by $f(T_{21}^{*})=0$ and<br><br>$G_{20}^{*}=\frac{(a_{20})^{(3)}G_{21}^{*}}{[(a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21}^{*})]}$ , $G_{22}^{*}=\frac{(a_{22})^{(3)}G_{21}^{*}}{[(a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21}^{*})]}$<br><br>$T_{20}^{*}=\frac{(b_{20})^{(3)}T_{21}^{*}}{[(b_{20}^{'})^{(3)}-(b_{20}^{''})^{(3)}(G_{23}{}^{*})]}$ , $T_{22}^{*}=\frac{(b_{22})^{(3)}T_{21}^{*}}{[(b_{22}^{'})^{(3)}-(b_{22}^{''})^{(3)}(G_{23}{}^{*})]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 515 |
| Finally we obtain the unique solution<br><br>$G_{25}^{*}$ given by $\varphi(G_{27})=0$ , $T_{25}^{*}$ given by $f(T_{25}^{*})=0$ and<br><br>$G_{24}^{*}=\frac{(a_{24})^{(4)}G_{25}^{*}}{[(a_{24}^{'})^{(4)}+(a_{24}^{''})^{(4)}(T_{25}^{*})]}$ , $G_{26}^{*}=\frac{(a_{26})^{(4)}G_{25}^{*}}{[(a_{26}^{'})^{(4)}+(a_{26}^{''})^{(4)}(T_{25}^{*})]}$ | 516 |
| $T_{24}^{*}=\frac{(b_{24})^{(4)}T_{25}^{*}}{[(b_{24}^{'})^{(4)}-(b_{24}^{''})^{(4)}((G_{27})^{*})]}$ , $T_{26}^{*}=\frac{(b_{26})^{(4)}T_{25}^{*}}{[(b_{26}^{'})^{(4)}-(b_{26}^{''})^{(4)}((G_{27})^{*})]}$ | 517 |

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| Obviously, these values represent an equilibrium solution of global equations | |
| Finally we obtain the unique solution<br><br>$G_{29}^{*}$ given by $\varphi((G_{31})^{*}) = 0$ , $T_{29}^{*}$ given by $f(T_{29}^{*}) = 0$ and<br><br>$G_{28}^{*} = \frac{(a_{28})^{(5)} G_{29}^{*}}{[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}^{*})]}$ , $G_{30}^{*} = \frac{(a_{30})^{(5)} G_{29}^{*}}{[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29}^{*})]}$ | 518 |
| $T_{28}^{*} = \frac{(b_{28})^{(5)} T_{29}^{*}}{[(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}((G_{31})^{*})]}$ , $T_{30}^{*} = \frac{(b_{30})^{(5)} T_{29}^{*}}{[(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}((G_{31})^{*})]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 519 |
| Finally we obtain the unique solution<br><br>$G_{33}^{*}$ given by $\varphi((G_{35})^{*}) = 0$ , $T_{33}^{*}$ given by $f(T_{33}^{*}) = 0$ and<br><br>$G_{32}^{*} = \frac{(a_{32})^{(6)} G_{33}^{*}}{[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33}^{*})]}$ , $G_{34}^{*} = \frac{(a_{34})^{(6)} G_{33}^{*}}{[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33}^{*})]}$ | 520 |
| $T_{32}^{*} = \frac{(b_{32})^{(6)} T_{33}^{*}}{[(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}((G_{35})^{*})]}$ , $T_{34}^{*} = \frac{(b_{34})^{(6)} T_{33}^{*}}{[(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}((G_{35})^{*})]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 521 |
| Finally we obtain the unique solution<br><br>$G_{37}^{*}$ given by $\varphi((G_{39})^{*}) = 0$ , $T_{37}^{*}$ given by $f(T_{37}^{*}) = 0$ and<br><br>$G_{36}^{*} = \frac{(a_{36})^{(7)} G_{37}^{*}}{[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37}^{*})]}$ , $G_{38}^{*} = \frac{(a_{38})^{(7)} G_{37}^{*}}{[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37}^{*})]}$<br><br>$T_{36}^{*} = \frac{(b_{36})^{(7)} T_{37}^{*}}{[(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}((G_{39})^{*})]}$ , $T_{38}^{*} = \frac{(b_{38})^{(7)} T_{37}^{*}}{[(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}((G_{39})^{*})]}$ | 522 |
| Finally we obtain the unique solution<br><br>$G_{41}^{*}$ given by $\varphi((G_{43})^{*}) = 0$ , $T_{41}^{*}$ given by $f(T_{41}^{*}) = 0$ and<br><br>$G_{40}^{*} = \frac{(a_{40})^{(8)} G_{41}^{*}}{[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41}^{*})]}$ , $G_{42}^{*} = \frac{(a_{42})^{(8)} G_{41}^{*}}{[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41}^{*})]}$<br><br>$T_{40}^{*} = \frac{(b_{40})^{(8)} T_{41}^{*}}{[(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}((G_{43})^{*})]}$ , $T_{42}^{*} = \frac{(b_{42})^{(8)} T_{41}^{*}}{[(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}((G_{43})^{*})]}$ | 523 |
| Finally we obtain the unique solution of 89 to 99 | 523 A |

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| $G_{45}^{*}$ given by $\varphi((G_{47})^{*}) = 0$ , $T_{45}^{*}$ given by $f(T_{45}^{*}) = 0$ and<br>$G_{44}^{*} = \frac{(a_{44})^{(9)} G_{45}^{*}}{[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45}^{*})]}$ , $G_{46}^{*} = \frac{(a_{46})^{(9)} G_{45}^{*}}{[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45}^{*})]}$<br>$T_{44}^{*} = \frac{(b_{44})^{(9)} T_{45}^{*}}{[(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}((G_{47})^{*})]}$ , $T_{46}^{*} = \frac{(b_{46})^{(9)} T_{45}^{*}}{[(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}((G_{47})^{*})]}$ | |
| **ASYMPTOTIC STABILITY ANALYSIS**<br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(1)}$ *and* $(b_i^{''})^{(1)}$ Belong to $C^{(1)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br>**Proof:** Denote<br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br>$G_i = G_i^{*} + \mathbb{G}_i$ , $T_i = T_i^{*} + \mathbb{T}_i$<br>$\frac{\partial (a_{14}^{''})^{(1)}}{\partial T_{14}}(T_{14}^{*}) = (q_{14})^{(1)}$ , $\frac{\partial (b_i^{''})^{(1)}}{\partial G_j}(G^{*}) = s_{ij}$ | 524 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{13}}{dt} = -\left((a_{13}^{'})^{(1)} + (p_{13})^{(1)}\right)\mathbb{G}_{13} + (a_{13})^{(1)}\mathbb{G}_{14} - (q_{13})^{(1)} G_{13}^{*}\mathbb{T}_{14}$ | 525 |
| $\frac{d\mathbb{G}_{14}}{dt} = -\left((a_{14}^{'})^{(1)} + (p_{14})^{(1)}\right)\mathbb{G}_{14} + (a_{14})^{(1)}\mathbb{G}_{13} - (q_{14})^{(1)} G_{14}^{*}\mathbb{T}_{14}$ | 526 |
| $\frac{d\mathbb{G}_{15}}{dt} = -\left((a_{15}^{'})^{(1)} + (p_{15})^{(1)}\right)\mathbb{G}_{15} + (a_{15})^{(1)}\mathbb{G}_{14} - (q_{15})^{(1)} G_{15}^{*}\mathbb{T}_{14}$ | 527 |
| $\frac{d\mathbb{T}_{13}}{dt} = -\left((b_{13}^{'})^{(1)} - (r_{13})^{(1)}\right)\mathbb{T}_{13} + (b_{13})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(13)(j)} T_{13}^{*}\mathbb{G}_j\right)$ | 528 |
| $\frac{d\mathbb{T}_{14}}{dt} = -\left((b_{14}^{'})^{(1)} - (r_{14})^{(1)}\right)\mathbb{T}_{14} + (b_{14})^{(1)}\mathbb{T}_{13} + \sum_{j=13}^{15}\left(s_{(14)(j)} T_{14}^{*}\mathbb{G}_j\right)$ | 529 |
| $\frac{d\mathbb{T}_{15}}{dt} = -\left((b_{15}^{'})^{(1)} - (r_{15})^{(1)}\right)\mathbb{T}_{15} + (b_{15})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(15)(j)} T_{15}^{*}\mathbb{G}_j\right)$ | 530 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}$ Belong to $C^{(2)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable | 531 |
| Proof: Denote<br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- | |
| $G_i = G_i^{*} + \mathbb{G}_i$ , $T_i = T_i^{*} + \mathbb{T}_i$ | 532 |

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| $\frac{\partial(a_{17}^{''})^{(2)}}{\partial \mathrm{T}_{17}}(\mathrm{T}_{17}^{*}) = (q_{17})^{(2)}$ , $\frac{\partial(b_i^{''})^{(2)}}{\partial \mathrm{G}_j}((G_{19})^{*}) = s_{ij}$ | 533 |
| taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{\mathrm{d}\mathbb{G}_{16}}{\mathrm{dt}} = -\left((a_{16}^{'})^{(2)} + (p_{16})^{(2)}\right)\mathbb{G}_{16} + (a_{16})^{(2)}\mathbb{G}_{17} - (q_{16})^{(2)}\mathrm{G}_{16}^{*}\mathbb{T}_{17}$ | 534 |
| $\frac{\mathrm{d}\mathbb{G}_{17}}{\mathrm{dt}} = -\left((a_{17}^{'})^{(2)} + (p_{17})^{(2)}\right)\mathbb{G}_{17} + (a_{17})^{(2)}\mathbb{G}_{16} - (q_{17})^{(2)}\mathrm{G}_{17}^{*}\mathbb{T}_{17}$ | 535 |
| $\frac{\mathrm{d}\mathbb{G}_{18}}{\mathrm{dt}} = -\left((a_{18}^{'})^{(2)} + (p_{18})^{(2)}\right)\mathbb{G}_{18} + (a_{18})^{(2)}\mathbb{G}_{17} - (q_{18})^{(2)}\mathrm{G}_{18}^{*}\mathbb{T}_{17}$ | 536 |
| $\frac{\mathrm{d}\mathbb{T}_{16}}{\mathrm{dt}} = -\left((b_{16}^{'})^{(2)} - (r_{16})^{(2)}\right)\mathbb{T}_{16} + (b_{16})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(16)(j)}\mathrm{T}_{16}^{*}\mathbb{G}_j\right)$ | 537 |
| $\frac{\mathrm{d}\mathbb{T}_{17}}{\mathrm{dt}} = -\left((b_{17}^{'})^{(2)} - (r_{17})^{(2)}\right)\mathbb{T}_{17} + (b_{17})^{(2)}\mathbb{T}_{16} + \sum_{j=16}^{18}\left(s_{(17)(j)}\mathrm{T}_{17}^{*}\mathbb{G}_j\right)$ | 538 |
| $\frac{\mathrm{d}\mathbb{T}_{18}}{\mathrm{dt}} = -\left((b_{18}^{'})^{(2)} - (r_{18})^{(2)}\right)\mathbb{T}_{18} + (b_{18})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(18)(j)}\mathrm{T}_{18}^{*}\mathbb{G}_j\right)$ | 539 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(3)}$ *and* $(b_i^{''})^{(3)}$ Belong to $C^{(3)}(\mathbb{R}_{+})$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^{*} + \mathbb{G}_i \quad , T_i = T_i^{*} + \mathbb{T}_i$<br><br>$\frac{\partial(a_{21}^{''})^{(3)}}{\partial T_{21}}(T_{21}^{*}) = (q_{21})^{(3)}$ , $\frac{\partial(b_i^{''})^{(3)}}{\partial G_j}((G_{23})^{*}) = s_{ij}$ | 540 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{20}}{dt} = -\left((a_{20}^{'})^{(3)} + (p_{20})^{(3)}\right)\mathbb{G}_{20} + (a_{20})^{(3)}\mathbb{G}_{21} - (q_{20})^{(3)}G_{20}^{*}\mathbb{T}_{21}$ | 541 |
| $\frac{d\mathbb{G}_{21}}{dt} = -\left((a_{21}^{'})^{(3)} + (p_{21})^{(3)}\right)\mathbb{G}_{21} + (a_{21})^{(3)}\mathbb{G}_{20} - (q_{21})^{(3)}G_{21}^{*}\mathbb{T}_{21}$ | 542 |
| $\frac{d\mathbb{G}_{22}}{dt} = -\left((a_{22}^{'})^{(3)} + (p_{22})^{(3)}\right)\mathbb{G}_{22} + (a_{22})^{(3)}\mathbb{G}_{21} - (q_{22})^{(3)}G_{22}^{*}\mathbb{T}_{21}$ | 543 |
| $\frac{d\mathbb{T}_{20}}{dt} = -\left((b_{20}^{'})^{(3)} - (r_{20})^{(3)}\right)\mathbb{T}_{20} + (b_{20})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(20)(j)}T_{20}^{*}\mathbb{G}_j\right)$ | 544 |
| $\frac{d\mathbb{T}_{21}}{dt} = -\left((b_{21}^{'})^{(3)} - (r_{21})^{(3)}\right)\mathbb{T}_{21} + (b_{21})^{(3)}\mathbb{T}_{20} + \sum_{j=20}^{22}\left(s_{(21)(j)}T_{21}^{*}\mathbb{G}_j\right)$ | 545 |
| $\frac{d\mathbb{T}_{22}}{dt} = -\left((b_{22}^{'})^{(3)} - (r_{22})^{(3)}\right)\mathbb{T}_{22} + (b_{22})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(22)(j)}T_{22}^{*}\mathbb{G}_j\right)$ | 546 |
| **ASYMPTOTIC STABILITY ANALYSIS** | 547 |

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| **Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(4)}$ $and$ $(b_i^{''})^{(4)}$ Belong to $C^{(4)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial (a_{25}^{''})^{(4)}}{\partial T_{25}}(T_{25}^*) = (q_{25})^{(4)}$ , $\frac{\partial (b_i^{''})^{(4)}}{\partial G_j}((G_{27})^*) = s_{ij}$ | 548 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{24}}{dt} = -\left((a_{24}^{'})^{(4)} + (p_{24})^{(4)}\right)\mathbb{G}_{24} + (a_{24})^{(4)}\mathbb{G}_{25} - (q_{24})^{(4)}G_{24}^*\mathbb{T}_{25}$ | 549 |
| $\frac{d\mathbb{G}_{25}}{dt} = -\left((a_{25}^{'})^{(4)} + (p_{25})^{(4)}\right)\mathbb{G}_{25} + (a_{25})^{(4)}\mathbb{G}_{24} - (q_{25})^{(4)}G_{25}^*\mathbb{T}_{25}$ | 550 |
| $\frac{d\mathbb{G}_{26}}{dt} = -\left((a_{26}^{'})^{(4)} + (p_{26})^{(4)}\right)\mathbb{G}_{26} + (a_{26})^{(4)}\mathbb{G}_{25} - (q_{26})^{(4)}G_{26}^*\mathbb{T}_{25}$ | 551 |
| $\frac{d\mathbb{T}_{24}}{dt} = -\left((b_{24}^{'})^{(4)} - (r_{24})^{(4)}\right)\mathbb{T}_{24} + (b_{24})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(24)(j)}T_{24}^*\mathbb{G}_j\right)$ | 552 |
| $\frac{d\mathbb{T}_{25}}{dt} = -\left((b_{25}^{'})^{(4)} - (r_{25})^{(4)}\right)\mathbb{T}_{25} + (b_{25})^{(4)}\mathbb{T}_{24} + \sum_{j=24}^{26}\left(s_{(25)(j)}T_{25}^*\mathbb{G}_j\right)$ | 553 |
| $\frac{d\mathbb{T}_{26}}{dt} = -\left((b_{26}^{'})^{(4)} - (r_{26})^{(4)}\right)\mathbb{T}_{26} + (b_{26})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(26)(j)}T_{26}^*\mathbb{G}_j\right)$ | 554 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 5:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(5)}$ $and$ $(b_i^{''})^{(5)}$ Belong to $C^{(5)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 555 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial (a_{29}^{''})^{(5)}}{\partial T_{29}}(T_{29}^*) = (q_{29})^{(5)}$ , $\frac{\partial (b_i^{''})^{(5)}}{\partial G_j}((G_{31})^*) = s_{ij}$ | 556 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{28}}{dt} = -\left((a_{28}^{'})^{(5)} + (p_{28})^{(5)}\right)\mathbb{G}_{28} + (a_{28})^{(5)}\mathbb{G}_{29} - (q_{28})^{(5)}G_{28}^*\mathbb{T}_{29}$ | 557 |
| $\frac{d\mathbb{G}_{29}}{dt} = -\left((a_{29}^{'})^{(5)} + (p_{29})^{(5)}\right)\mathbb{G}_{29} + (a_{29})^{(5)}\mathbb{G}_{28} - (q_{29})^{(5)}G_{29}^*\mathbb{T}_{29}$ | 558 |
| $\frac{d\mathbb{G}_{30}}{dt} = -\left((a_{30}^{'})^{(5)} + (p_{30})^{(5)}\right)\mathbb{G}_{30} + (a_{30})^{(5)}\mathbb{G}_{29} - (q_{30})^{(5)}G_{30}^*\mathbb{T}_{29}$ | 559 |

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| $\frac{d\mathbb{T}_{28}}{dt} = -\left((b_{28}^{'})^{(5)} - (r_{28})^{(5)}\right)\mathbb{T}_{28} + (b_{28})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(28)(j)}T_{28}^{*}\mathbb{G}_{j}\right)$ | 560 |
| $\frac{d\mathbb{T}_{29}}{dt} = -\left((b_{29}^{'})^{(5)} - (r_{29})^{(5)}\right)\mathbb{T}_{29} + (b_{29})^{(5)}\mathbb{T}_{28} + \sum_{j=28}^{30}\left(s_{(29)(j)}T_{29}^{*}\mathbb{G}_{j}\right)$ | 561 |
| $\frac{d\mathbb{T}_{30}}{dt} = -\left((b_{30}^{'})^{(5)} - (r_{30})^{(5)}\right)\mathbb{T}_{30} + (b_{30})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(30)(j)}T_{30}^{*}\mathbb{G}_{j}\right)$ | 562 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 6:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(6)}$ $and$ $(b_i^{''})^{(6)}$ Belong to $C^{(6)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 563 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial(a_{33}^{''})^{(6)}}{\partial T_{33}}(T_{33}^{*}) = (q_{33})^{(6)}$ , $\frac{\partial(b_i^{''})^{(6)}}{\partial G_j}((G_{35})^{*}) = s_{ij}$ | 564 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{32}}{dt} = -\left((a_{32}^{'})^{(6)} + (p_{32})^{(6)}\right)\mathbb{G}_{32} + (a_{32})^{(6)}\mathbb{G}_{33} - (q_{32})^{(6)}G_{32}^{*}\mathbb{T}_{33}$ | 565 |
| $\frac{d\mathbb{G}_{33}}{dt} = -\left((a_{33}^{'})^{(6)} + (p_{33})^{(6)}\right)\mathbb{G}_{33} + (a_{33})^{(6)}\mathbb{G}_{32} - (q_{33})^{(6)}G_{33}^{*}\mathbb{T}_{33}$ | 566 |
| $\frac{d\mathbb{G}_{34}}{dt} = -\left((a_{34}^{'})^{(6)} + (p_{34})^{(6)}\right)\mathbb{G}_{34} + (a_{34})^{(6)}\mathbb{G}_{33} - (q_{34})^{(6)}G_{34}^{*}\mathbb{T}_{33}$ | 567 |
| $\frac{d\mathbb{T}_{32}}{dt} = -\left((b_{32}^{'})^{(6)} - (r_{32})^{(6)}\right)\mathbb{T}_{32} + (b_{32})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(32)(j)}T_{32}^{*}\mathbb{G}_{j}\right)$ | 568 |
| $\frac{d\mathbb{T}_{33}}{dt} = -\left((b_{33}^{'})^{(6)} - (r_{33})^{(6)}\right)\mathbb{T}_{33} + (b_{33})^{(6)}\mathbb{T}_{32} + \sum_{j=32}^{34}\left(s_{(33)(j)}T_{33}^{*}\mathbb{G}_{j}\right)$ | 569 |
| $\frac{d\mathbb{T}_{34}}{dt} = -\left((b_{34}^{'})^{(6)} - (r_{34})^{(6)}\right)\mathbb{T}_{34} + (b_{34})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(34)(j)}T_{34}^{*}\mathbb{G}_{j}\right)$ | 570 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 7:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(7)}$ $and$ $(b_i^{''})^{(7)}$ Belong to $C^{(7)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 571 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$ | 572 |

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| $\frac{\partial (a_{37}^{''})^{(7)}}{\partial T_{37}}(T_{37}^{*}) = (q_{37})^{(7)}$ , $\frac{\partial (b_i^{''})^{(7)}}{\partial G_j}\left((G_{39})^{**}\right) = s_{ij}$ | |
| Then taking into account equations and neglecting the terms of power 2, we obtain from | |
| $\frac{d\mathbb{G}_{36}}{dt} = -\left((a_{36}^{'})^{(7)} + (p_{36})^{(7)}\right)\mathbb{G}_{36} + (a_{36})^{(7)}\mathbb{G}_{37} - (q_{36})^{(7)}G_{36}^{*}\mathbb{T}_{37}$ | 573 |
| $\frac{d\mathbb{G}_{37}}{dt} = -\left((a_{37}^{'})^{(7)} + (p_{37})^{(7)}\right)\mathbb{G}_{37} + (a_{37})^{(7)}\mathbb{G}_{36} - (q_{37})^{(7)}G_{37}^{*}\mathbb{T}_{37}$ | 574 |
| $\frac{d\mathbb{G}_{38}}{dt} = -\left((a_{38}^{'})^{(7)} + (p_{38})^{(7)}\right)\mathbb{G}_{38} + (a_{38})^{(7)}\mathbb{G}_{37} - (q_{38})^{(7)}G_{38}^{*}\mathbb{T}_{37}$ | 575 |
| $\frac{d\mathbb{T}_{36}}{dt} = -\left((b_{36}^{'})^{(7)} - (r_{36})^{(7)}\right)\mathbb{T}_{36} + (b_{36})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(36)(j)}T_{36}^{*}\mathbb{G}_j\right)$ | 576 |
| $\frac{d\mathbb{T}_{37}}{dt} = -\left((b_{37}^{'})^{(7)} - (r_{37})^{(7)}\right)\mathbb{T}_{37} + (b_{37})^{(7)}\mathbb{T}_{36} + \sum_{j=36}^{38}\left(s_{(37)(j)}T_{37}^{*}\mathbb{G}_j\right)$ | 578 |
| $\frac{d\mathbb{T}_{38}}{dt} = -\left((b_{38}^{'})^{(7)} - (r_{38})^{(7)}\right)\mathbb{T}_{38} + (b_{38})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(38)(j)}T_{38}^{*}\mathbb{G}_j\right)$ | 579 |
| Obviously, these values represent an equilibrium solution<br><br>**ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 8:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(8)}$ *and* $(b_i^{''})^{(8)}$ Belong to $C^{(8)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^{*} + \mathbb{G}_i \quad , T_i = T_i^{*} + \mathbb{T}_i$<br><br>$\frac{\partial (a_{41}^{''})^{(8)}}{\partial T_{41}}(T_{41}^{*}) = (q_{41})^{(8)}$ , $\frac{\partial (b_i^{''})^{(8)}}{\partial G_j}\left((G_{43})^{*}\right) = s_{ij}$ | 580 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{40}}{dt} = -\left((a_{40}^{'})^{(8)} + (p_{40})^{(8)}\right)\mathbb{G}_{40} + (a_{40})^{(8)}\mathbb{G}_{41} - (q_{40})^{(8)}G_{40}^{*}\mathbb{T}_{41}$ | 581 |
| $\frac{d\mathbb{G}_{41}}{dt} = -\left((a_{41}^{'})^{(8)} + (p_{41})^{(8)}\right)\mathbb{G}_{41} + (a_{41})^{(8)}\mathbb{G}_{40} - (q_{41})^{(8)}G_{41}^{*}\mathbb{T}_{41}$ | 582 |
| $\frac{d\mathbb{G}_{42}}{dt} = -\left((a_{42}^{'})^{(8)} + (p_{42})^{(8)}\right)\mathbb{G}_{42} + (a_{42})^{(8)}\mathbb{G}_{41} - (q_{42})^{(8)}G_{42}^{*}\mathbb{T}_{41}$ | 583 |
| $\frac{d\mathbb{T}_{40}}{dt} = -\left((b_{40}^{'})^{(8)} - (r_{40})^{(8)}\right)\mathbb{T}_{40} + (b_{40})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(40)(j)}T_{40}^{*}\mathbb{G}_j\right)$ | 584 |

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|---|---|
| $\frac{d\mathbb{T}_{41}}{dt} = -\left((b'_{41})^{(8)} - (r_{41})^{(8)}\right)\mathbb{T}_{41} + (b_{41})^{(8)}\mathbb{T}_{40} + \sum_{j=40}^{42}\left(s_{(41)(j)}T^*_{41}\mathbb{G}_j\right)$ | 585 |
| $\frac{d\mathbb{T}_{42}}{dt} = -\left((b'_{42})^{(8)} - (r_{42})^{(8)}\right)\mathbb{T}_{42} + (b_{42})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(42)(j)}T^*_{42}\mathbb{G}_j\right)$ | 586 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 9:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(9)}$ *and* $(b_i^{''})^{(9)}$ Belong to $C^{(9)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br>**Proof: Denote** | 586 A |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial(a_{45}^{''})^{(9)}}{\partial T_{45}}(T^*_{45}) = (q_{45})^{(9)} \;,\; \frac{\partial(b_i^{''})^{(9)}}{\partial G_j}((G_{47})^*) = s_{ij}$ | |
| Then taking into account equations 89 to 99 and neglecting the terms of power 2, we obtain from 99 to 44 | |
| $\frac{d\mathbb{G}_{44}}{dt} = -\left((a'_{44})^{(9)} + (p_{44})^{(9)}\right)\mathbb{G}_{44} + (a_{44})^{(9)}\mathbb{G}_{45} - (q_{44})^{(9)}G^*_{44}\mathbb{T}_{45}$ | 586 B |
| $\frac{d\mathbb{G}_{45}}{dt} = -\left((a'_{45})^{(9)} + (p_{45})^{(9)}\right)\mathbb{G}_{45} + (a_{45})^{(9)}\mathbb{G}_{44} - (q_{45})^{(9)}G^*_{45}\mathbb{T}_{45}$ | 586 C |
| $\frac{d\mathbb{G}_{46}}{dt} = -\left((a'_{46})^{(9)} + (p_{46})^{(9)}\right)\mathbb{G}_{46} + (a_{46})^{(9)}\mathbb{G}_{45} - (q_{46})^{(9)}G^*_{46}\mathbb{T}_{45}$ | 586 D |
| $\frac{d\mathbb{T}_{44}}{dt} = -\left((b'_{44})^{(9)} - (r_{44})^{(9)}\right)\mathbb{T}_{44} + (b_{44})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(44)(j)}T^*_{44}\mathbb{G}_j\right)$ | 586 E |
| $\frac{d\mathbb{T}_{45}}{dt} = -\left((b'_{45})^{(9)} - (r_{45})^{(9)}\right)\mathbb{T}_{45} + (b_{45})^{(9)}\mathbb{T}_{44} + \sum_{j=44}^{46}\left(s_{(45)(j)}T^*_{45}\mathbb{G}_j\right)$ | 586 F |
| $\frac{d\mathbb{T}_{46}}{dt} = -\left((b'_{46})^{(9)} - (r_{46})^{(9)}\right)\mathbb{T}_{46} + (b_{46})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(46)(j)}T^*_{46}\mathbb{G}_j\right)$ | 586 G |
| **The characteristic equation of this system is** | 587 |
| $\left((\lambda)^{(1)} + (b'_{15})^{(1)} - (r_{15})^{(1)}\right)\{\left((\lambda)^{(1)} + (a'_{15})^{(1)} + (p_{15})^{(1)}\right)$<br><br>$\left[\left((\lambda)^{(1)} + (a'_{13})^{(1)} + (p_{13})^{(1)}\right)(q_{14})^{(1)}G^*_{14} + (a_{14})^{(1)}(q_{13})^{(1)}G^*_{13}\right)\Big]$<br><br>$\left((\lambda)^{(1)} + (b'_{13})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(14)}T^*_{14} + (b_{14})^{(1)}s_{(13),(14)}T^*_{14}\Big)$<br><br>$+\left((\lambda)^{(1)} + (a'_{14})^{(1)} + (p_{14})^{(1)}\right)(q_{13})^{(1)}G^*_{13} + (a_{13})^{(1)}(q_{14})^{(1)}G^*_{14}\Big)$<br><br>$\left((\lambda)^{(1)} + (b'_{13})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(13)}T^*_{14} + (b_{14})^{(1)}s_{(13),(13)}T^*_{13}\Big)$<br><br>$\left((\lambda)^{(1)}\right)^2 + \left((a'_{13})^{(1)} + (a'_{14})^{(1)} + (p_{13})^{(1)} + (p_{14})^{(1)}\right)(\lambda)^{(1)}\Big)$ | |

$$\left((\lambda)^{(1)}\right)^2+\left((b_{13}')^{(1)}+(b_{14}')^{(1)}-(r_{13})^{(1)}+(r_{14})^{(1)}\right)(\lambda)^{(1)}\Big)$$

$$+\left(\left((\lambda)^{(1)}\right)^2+\left((a_{13}')^{(1)}+(a_{14}')^{(1)}+(p_{13})^{(1)}+(p_{14})^{(1)}\right)(\lambda)^{(1)}\right)(q_{15})^{(1)}G_{15}$$

$$+\left((\lambda)^{(1)}+(a_{13}')^{(1)}+(p_{13})^{(1)}\right)\left((a_{15})^{(1)}(q_{14})^{(1)}G_{14}^{*}+(a_{14})^{(1)}(a_{15})^{(1)}(q_{13})^{(1)}G_{13}^{*}\right)$$

$$\left((\lambda)^{(1)}+(b_{13}')^{(1)}-(r_{13})^{(1)}\right)s_{(14),(15)}T_{14}^{*}+(b_{14})^{(1)}s_{(13),(15)}T_{13}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(2)}+(b_{18}')^{(2)}-(r_{18})^{(2)}\}\{\left((\lambda)^{(2)}+(a_{18}')^{(2)}+(p_{18})^{(2)}\right)$$

$$\left[\left(\left((\lambda)^{(2)}+(a_{16}')^{(2)}+(p_{16})^{(2)}\right)(q_{17})^{(2)}\mathrm{G}_{17}^{*}+(a_{17})^{(2)}(q_{16})^{(2)}\mathrm{G}_{16}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(2)}+(b_{16}')^{(2)}-(r_{16})^{(2)}\right)s_{(17),(17)}\mathrm{T}_{17}^{*}+(b_{17})^{(2)}s_{(16),(17)}\mathrm{T}_{17}^{*}\right)$$

$$+\left(\left((\lambda)^{(2)}+(a_{17}')^{(2)}+(p_{17})^{(2)}\right)(q_{16})^{(2)}\mathrm{G}_{16}^{*}+(a_{16})^{(2)}(q_{17})^{(2)}\mathrm{G}_{17}^{*}\right)$$

$$\left(\left((\lambda)^{(2)}+(b_{16}')^{(2)}-(r_{16})^{(2)}\right)s_{(17),(16)}\mathrm{T}_{17}^{*}+(b_{17})^{(2)}s_{(16),(16)}\mathrm{T}_{16}^{*}\right)$$

$$\left(\left((\lambda)^{(2)}\right)^2+\left((a_{16}')^{(2)}+(a_{17}')^{(2)}+(p_{16})^{(2)}+(p_{17})^{(2)}\right)(\lambda)^{(2)}\right)$$

$$\left(\left((\lambda)^{(2)}\right)^2+\left((b_{16}')^{(2)}+(b_{17}')^{(2)}-(r_{16})^{(2)}+(r_{17})^{(2)}\right)(\lambda)^{(2)}\right)$$

$$+\left(\left((\lambda)^{(2)}\right)^2+\left((a_{16}')^{(2)}+(a_{17}')^{(2)}+(p_{16})^{(2)}+(p_{17})^{(2)}\right)(\lambda)^{(2)}\right)(q_{18})^{(2)}\mathrm{G}_{18}$$

$$+\left((\lambda)^{(2)}+(a_{16}')^{(2)}+(p_{16})^{(2)}\right)\left((a_{18})^{(2)}(q_{17})^{(2)}\mathrm{G}_{17}^{*}+(a_{17})^{(2)}(a_{18})^{(2)}(q_{16})^{(2)}\mathrm{G}_{16}^{*}\right)$$

$$\left((\lambda)^{(2)}+(b_{16}')^{(2)}-(r_{16})^{(2)}\right)s_{(17),(18)}\mathrm{T}_{17}^{*}+(b_{17})^{(2)}s_{(16),(18)}\mathrm{T}_{16}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(3)}+(b_{22}')^{(3)}-(r_{22})^{(3)}\}\{\left((\lambda)^{(3)}+(a_{22}')^{(3)}+(p_{22})^{(3)}\right)$$

$$\left[\left(\left((\lambda)^{(3)}+(a_{20}')^{(3)}+(p_{20})^{(3)}\right)(q_{21})^{(3)}G_{21}^{*}+(a_{21})^{(3)}(q_{20})^{(3)}G_{20}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(21)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(21)}T_{21}^{*}\right)$$

$$+\left(\left((\lambda)^{(3)}+(a_{21}')^{(3)}+(p_{21})^{(3)}\right)(q_{20})^{(3)}G_{20}^{*}+(a_{20})^{(3)}(q_{21})^{(1)}G_{21}^{*}\right)$$

$$\left(\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(20)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(20)}T_{20}^{*}\right)$$

$$\left(\left((\lambda)^{(3)}\right)^2+\left((a_{20}')^{(3)}+(a_{21}')^{(3)}+(p_{20})^{(3)}+(p_{21})^{(3)}\right)(\lambda)^{(3)}\right)$$

$$\left(\left((\lambda)^{(3)}\right)^2+\left((b_{20}')^{(3)}+(b_{21}')^{(3)}-(r_{20})^{(3)}+(r_{21})^{(3)}\right)(\lambda)^{(3)}\right)$$

$$+\left(\left((\lambda)^{(3)}\right)^2+\left((a_{20}')^{(3)}+(a_{21}')^{(3)}+(p_{20})^{(3)}+(p_{21})^{(3)}\right)(\lambda)^{(3)}\right)(q_{22})^{(3)}G_{22}$$

$$+\left((\lambda)^{(3)}+(a_{20}')^{(3)}+(p_{20})^{(3)}\right)\left((a_{22})^{(3)}(q_{21})^{(3)}G_{21}^{*}+(a_{21})^{(3)}(a_{22})^{(3)}(q_{20})^{(3)}G_{20}^{*}\right)$$

$$\left(\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(22)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(22)}T_{20}^{*}\right)\}=0$$

$$+$$

$$(\lambda)^{(4)}+(b_{26}')^{(4)}-(r_{26})^{(4)}\{\left((\lambda)^{(4)}+(a_{26}')^{(4)}+(p_{26})^{(4)}\right)$$

$$\left[\left(\left((\lambda)^{(4)}+(a_{24}')^{(4)}+(p_{24})^{(4)}\right)(q_{25})^{(4)}G_{25}^{*}+(a_{25})^{(4)}(q_{24})^{(4)}G_{24}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(4)}+(b_{24}')^{(4)}-(r_{24})^{(4)}\right)s_{(25),(25)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(25)}T_{25}^{*}\right)$$

$$+\left(\left((\lambda)^{(4)}+(a_{25}')^{(4)}+(p_{25})^{(4)}\right)(q_{24})^{(4)}G_{24}^{*}+(a_{24})^{(4)}(q_{25})^{(4)}G_{25}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}+(b_{24}')^{(4)}-(r_{24})^{(4)}\right)s_{(25),(24)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(24)}T_{24}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}\right)^2+\left((a_{24}')^{(4)}+(a_{25}')^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)$$

$$\left(\left((\lambda)^{(4)}\right)^2+\left((b_{24}')^{(4)}+(b_{25}')^{(4)}-(r_{24})^{(4)}+(r_{25})^{(4)}\right)(\lambda)^{(4)}\right)$$

$$+\left(\left((\lambda)^{(4)}\right)^2+\left((a_{24}')^{(4)}+(a_{25}')^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)(q_{26})^{(4)}G_{26}$$

$$+\left((\lambda)^{(4)}+(a_{24}')^{(4)}+(p_{24})^{(4)}\right)\left((a_{26})^{(4)}(q_{25})^{(4)}G_{25}^{*}+(a_{25})^{(4)}(a_{26})^{(4)}(q_{24})^{(4)}G_{24}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}+(b_{24}')^{(4)}-(r_{24})^{(4)}\right)s_{(25),(26)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(26)}T_{24}^{*}\right)\}=0$$

$$+$$

$$(\lambda)^{(5)}+(b_{30}')^{(5)}-(r_{30})^{(5)}\{\left((\lambda)^{(5)}+(a_{30}')^{(5)}+(p_{30})^{(5)}\right)$$

$$\left[\left(\left((\lambda)^{(5)}+(a_{28}')^{(5)}+(p_{28})^{(5)}\right)(q_{29})^{(5)}G_{29}^{*}+(a_{29})^{(5)}(q_{28})^{(5)}G_{28}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(5)}+(b_{28}')^{(5)}-(r_{28})^{(5)}\right)s_{(29),(29)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(29)}T_{29}^{*}\right)$$

$$+\left(\left((\lambda)^{(5)}+(a_{29}')^{(5)}+(p_{29})^{(5)}\right)(q_{28})^{(5)}G_{28}^{*}+(a_{28})^{(5)}(q_{29})^{(5)}G_{29}^{*}\right)$$

$$\left(\left((\lambda)^{(5)}+(b_{28}')^{(5)}-(r_{28})^{(5)}\right)s_{(29),(28)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(28)}T_{28}^{*}\right)$$

$$\left(\left((\lambda)^{(5)}\right)^2+\left((a_{28}')^{(5)}+(a_{29}')^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)$$

$$\left((\lambda)^{(5)}\right)^2+\left((b_{28}')^{(5)}+(b_{29}')^{(5)}-(r_{28})^{(5)}+(r_{29})^{(5)}\right)(\lambda)^{(5)}\Big)$$

$$+\left(\left((\lambda)^{(5)}\right)^2+\left((a_{28}')^{(5)}+(a_{29}')^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)(q_{30})^{(5)}G_{30}$$

$$+\left((\lambda)^{(5)}+(a_{28}')^{(5)}+(p_{28})^{(5)}\right)\left((a_{30})^{(5)}(q_{29})^{(5)}G_{29}^{*}+(a_{29})^{(5)}(a_{30})^{(5)}(q_{28})^{(5)}G_{28}^{*}\right)$$

$$\left((\lambda)^{(5)}+(b_{28}')^{(5)}-(r_{28})^{(5)}\right)s_{(29),(30)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(30)}T_{28}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(6)}+(b_{34}')^{(6)}-(r_{34})^{(6)}\{\left((\lambda)^{(6)}+(a_{34}')^{(6)}+(p_{34})^{(6)}\right)$$

$$\left[\left((\lambda)^{(6)}+(a_{32}')^{(6)}+(p_{32})^{(6)}\right)(q_{33})^{(6)}G_{33}^{*}+(a_{33})^{(6)}(q_{32})^{(6)}G_{32}^{*}\right)\Big]$$

$$\left((\lambda)^{(6)}+(b_{32}')^{(6)}-(r_{32})^{(6)}\right)s_{(33),(33)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(33)}T_{33}^{*}\Big)$$

$$+\left((\lambda)^{(6)}+(a_{33}')^{(6)}+(p_{33})^{(6)}\right)(q_{32})^{(6)}G_{32}^{*}+(a_{32})^{(6)}(q_{33})^{(6)}G_{33}^{*}\Big)$$

$$\left((\lambda)^{(6)}+(b_{32}')^{(6)}-(r_{32})^{(6)}\right)s_{(33),(32)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(32)}T_{32}^{*}\Big)$$

$$\left((\lambda)^{(6)}\right)^2+\left((a_{32}')^{(6)}+(a_{33}')^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\Big)$$

$$\left((\lambda)^{(6)}\right)^2+\left((b_{32}')^{(6)}+(b_{33}')^{(6)}-(r_{32})^{(6)}+(r_{33})^{(6)}\right)(\lambda)^{(6)}\Big)$$

$$+\left(\left((\lambda)^{(6)}\right)^2+\left((a_{32}')^{(6)}+(a_{33}')^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\right)(q_{34})^{(6)}G_{34}$$

$$+\left((\lambda)^{(6)}+(a_{32}')^{(6)}+(p_{32})^{(6)}\right)\left((a_{34})^{(6)}(q_{33})^{(6)}G_{33}^{*}+(a_{33})^{(6)}(a_{34})^{(6)}(q_{32})^{(6)}G_{32}^{*}\right)$$

$$\left((\lambda)^{(6)}+(b_{32}')^{(6)}-(r_{32})^{(6)}\right)s_{(33),(34)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(34)}T_{32}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(7)}+(b_{38}')^{(7)}-(r_{38})^{(7)}\{\left((\lambda)^{(7)}+(a_{38}')^{(7)}+(p_{38})^{(7)}\right)$$

$$\left[\left((\lambda)^{(7)}+(a_{36}')^{(7)}+(p_{36})^{(7)}\right)(q_{37})^{(7)}G_{37}^{*}+(a_{37})^{(7)}(q_{36})^{(7)}G_{36}^{*}\right)\Big]$$

$$\left((\lambda)^{(7)}+(b_{36}')^{(7)}-(r_{36})^{(7)}\right)s_{(37),(37)}T_{37}^{*}+(b_{37})^{(7)}s_{(36),(37)}T_{37}^{*}\Big)$$

$$+\left((\lambda)^{(7)}+(a_{37}')^{(7)}+(p_{37})^{(7)}\right)(q_{36})^{(7)}G_{36}^{*}+(a_{36})^{(7)}(q_{37})^{(7)}G_{37}^{*}\Big)$$

$$\left((\lambda)^{(7)}+(b_{36}')^{(7)}-(r_{36})^{(7)}\right)s_{(37),(36)}T_{37}^{*}+(b_{37})^{(7)}s_{(36),(36)}T_{36}^{*}\Big)$$

$$\left((\lambda)^{(7)}\right)^2+\left((a_{36}')^{(7)}+(a_{37}')^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\Big)$$

$$\left((\lambda)^{(7)}\right)^2+\left((b'_{36})^{(7)}+(b'_{37})^{(7)}-(r_{36})^{(7)}+(r_{37})^{(7)}\right)(\lambda)^{(7)}\Big)$$

$$+\left(\left((\lambda)^{(7)}\right)^2+\left((a'_{36})^{(7)}+(a'_{37})^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\right)(q_{38})^{(7)}G_{38}$$

$$+\left((\lambda)^{(7)}+(a'_{36})^{(7)}+(p_{36})^{(7)}\right)\left((a_{38})^{(7)}(q_{37})^{(7)}G^*_{37}+(a_{37})^{(7)}(a_{38})^{(7)}(q_{36})^{(7)}G^*_{36}\right)$$

$$\left((\lambda)^{(7)}+(b'_{36})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(38)}T^*_{37}+(b_{37})^{(7)}s_{(36),(38)}T^*_{36}\Big)\}=0$$

$$+$$

$$(\lambda)^{(8)}+(b'_{42})^{(8)}-(r_{42})^{(8)}\}\{\left((\lambda)^{(8)}+(a'_{42})^{(8)}+(p_{42})^{(8)}\right)$$

$$\left[\left(\left((\lambda)^{(8)}+(a'_{40})^{(8)}+(p_{40})^{(8)}\right)(q_{41})^{(8)}G^*_{41}+(a_{41})^{(8)}(q_{40})^{(8)}G^*_{40}\right)\right]$$

$$\left(\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(41)}T^*_{41}+(b_{41})^{(8)}s_{(40),(41)}T^*_{41}\right)$$

$$+\left(\left((\lambda)^{(8)}+(a'_{41})^{(8)}+(p_{41})^{(8)}\right)(q_{40})^{(8)}G^*_{40}+(a_{40})^{(8)}(q_{41})^{(8)}G^*_{41}\right)$$

$$\left(\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(40)}T^*_{41}+(b_{41})^{(8)}s_{(40),(40)}T^*_{40}\right)$$

$$\left((\lambda)^{(8)}\right)^2+\left((a'_{40})^{(8)}+(a'_{41})^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\Big)$$

$$\left((\lambda)^{(8)}\right)^2+\left((b'_{40})^{(8)}+(b'_{41})^{(8)}-(r_{40})^{(8)}+(r_{41})^{(8)}\right)(\lambda)^{(8)}\Big)$$

$$+\left(\left((\lambda)^{(8)}\right)^2+\left((a'_{40})^{(8)}+(a'_{41})^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\right)(q_{42})^{(8)}G_{42}$$

$$+\left((\lambda)^{(8)}+(a'_{40})^{(8)}+(p_{40})^{(8)}\right)\left((a_{42})^{(8)}(q_{41})^{(8)}G^*_{41}+(a_{41})^{(8)}(a_{42})^{(8)}(q_{40})^{(8)}G^*_{40}\right)$$

$$\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(42)}T^*_{41}+(b_{41})^{(8)}s_{(40),(42)}T^*_{40}\Big)\}=0$$

$$+$$

$$(\lambda)^{(9)}+(b'_{46})^{(9)}-(r_{46})^{(9)}\}\{\left((\lambda)^{(9)}+(a'_{46})^{(9)}+(p_{46})^{(9)}\right)$$

$$\left[\left(\left((\lambda)^{(9)}+(a'_{44})^{(9)}+(p_{44})^{(9)}\right)(q_{45})^{(9)}G^*_{45}+(a_{45})^{(9)}(q_{44})^{(9)}G^*_{44}\right)\right]$$

$$\left(\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(45)}T^*_{45}+(b_{45})^{(9)}s_{(44),(45)}T^*_{45}\right)$$

$$+\left(\left((\lambda)^{(9)}+(a'_{45})^{(9)}+(p_{45})^{(9)}\right)(q_{44})^{(9)}G^*_{44}+(a_{44})^{(9)}(q_{45})^{(9)}G^*_{45}\right)$$

$$\left(\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(44)}T^*_{45}+(b_{45})^{(9)}s_{(44),(44)}T^*_{44}\right)$$

$$\left((\lambda)^{(9)}\right)^2+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\Big)$$

$$\left((\lambda)^{(9)}\right)^2+\left((b'_{44})^{(9)}+(b'_{45})^{(9)}-(r_{44})^{(9)}+(r_{45})^{(9)}\right)(\lambda)^{(9)}\Big)$$

$$+\left(\left((\lambda)^{(9)}\right)^2+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\right)(q_{46})^{(9)}G_{46}$$

$$+\left((\lambda)^{(9)}+(a'_{44})^{(9)}+(p_{44})^{(9)}\right)\left((a_{46})^{(9)}(q_{45})^{(9)}G^*_{45}+(a_{45})^{(9)}(a_{46})^{(9)}(q_{44})^{(9)}G^*_{44}\right)$$

$$\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(46)}T^*_{45}+(b_{45})^{(9)}s_{(44),(46)}T^*_{44}\Big)\}=0$$

**And as one sees, all the coefficients are positive. It follows that all the roots have negative real part, and this proves the theorem.**

Note: **Same dovetailing explanation holds good for d/dt (partial differential with respect to t), d/dt,d^2/dt^2 (acceralation: double dot). Such an exposition is helpful in optimisation problems, duality of motion, wave motion a homogeneous functions studies, and control theory which we intend to incorporate in future.** It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## Persistence, Coexistence And Thrusty Permanence For Interacting Structured Particles, Transactions In Finance, Econometric Productivity Functions, Environment, Interfaces, Cosmological Interactions Et Al

## Josef Hofbauer And Sebastian J. Schreiber Model Thereof: A Functional Appraisal

ABSTRACT. **JOSEF HOFBAUER AND SEBASTIAN J. SCHREIBER** studied the dynamics of interacting structured systems (in their case population dynamics) can be modeled by $\frac{dxi}{dt}= A_i (x) xi$ where $x_i \epsilon R^n_{i,} x = (x_1,\ldots,x_k)$, and $A_i(x)$ are matrices with non-negative off- diagonal entries. These models are permanent if there exists

a positive global attractor and are robustly permanent if they remain permanent following perturbations of $A_i\ (x)$. Necessary and sufficient conditions for robust permanence are derived using dominant Lyapunov exponents $X_i\ (\mu)$ of the $A_i\ (x)$ with respect to invariant measures $\mu$. The necessary conditions requires $\max_i\ X_i(\mu)>0$ for all ergodic measures with support in the boundary of the non-negative cone. The sufficient condition requires that the boundary admits a Morse decomposition such that $\max_i$ Xi$(\mu)>0$ for all invariant measures $\mu$ supported by a component of the Morse decomposition. When the Morse components are Axiom A, uniquely ergodic, or support all but one population, the necessary and sufficient conditions are equivalent. Applications to spatial ecology, epidemiology, and gene networks are given. Appeared in *Journal of Differential Equations*, 248, 1955-1971 (2010) It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## INTRODUCTION

A fundamental issue in population biology is what are the minimal conditions to ensure the long-term survivorship for interacting populations whether they be viral particles, bio-chemicals, plants, or animals. When these conditions are met the interacting populations are said to persist or coexist. Since the pioneering work of Lotka and Volterra on competitive and predator prey interactions, Thompson, Nicholson, and Bailey on host-parasite interactions, and Kermack and McKendrick on disease outbreaks, nonlinear difference and differential equations have been used to understand conditions for population persistence [1-5]. One particularly important form of persistence for deterministic models is *permanence* or *uniform persistence* which corresponds to the existence of a global attractor bounded away from extinction of one or more species. When such an attractor exists, interacting populations are able to recover from "vigorous shake ups" of the of the population state. Permanence has been characterized from a topological perspective [6-8] and with average Lyapunov functions [9-12] for differential equation models of the form

$$\frac{dxi}{dt} = x_i f_i(x) \qquad i = 1,\ldots,k \qquad (6)$$

Where $x_i$ is the abundance of population $i$, $x = (x_1,\ldots,x_k)$, and $f_i(x)$ is the per-capita growth of population $i$.

An important extension of the concept of the permanence is *robust permanence* [13-16] which requires that (1) remain permanent after sufficiently small perturbations of the per-capita growth rates $f_i$. The importance of robust permanence as one crude form of structural stability for population models are approximations to reality. Consequently, if nearby models (e.g. more realistic models) are not permanent despite the focal model being permanent, and then one can draw few conclusions about the persistence of the biological system being approximated by the model. One can view robust permanence as one crude form of structural stability for debits-credits particles, propositions, transactions in the celestial subterranean realm and ceratoid dualism, finance like debiting say a current account and crediting to another account, or recrediting to the account which had been earlier debited, further conduction of the cash credit account, black hole formation, guzzling of inventories for the output, creation of income from output, reinvestment of income to investment,gormadizing of matter by black holes, and emission of radiation therefore, gamma ray bursts, in fact most of the transactions and interactions in the universe are conservative and quantum mechanical universe, something like a account acting as predator and prey. More often than not as we have seen from Haimovici model such models could be applied to the problems, propositions, which are not subsumed by the concomitant problems in the bearing. the population models. For this perspective, it is not unexpected that there are permanent systems that cannot be approximated by robustly permanent systems [17]. In [14, 16], criteria for robust permanence were developed with respect to the average per-capita growth rates $\int f_i(x)d\mu(x)$ with respect to invariant probability measures $\mu$ supported on the boundary of the positive orthant (i.e. where one or more Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates are extinct). Roughly, these criteria for robust permanence require that the average per-capita growth rate (i.e. invasion rate) is positive for some missing species.

While (1) can account for many types of Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates. Interactions, it assumes that all individual within populations are exactly the same. However, theoretical biologists have long recognized that different individuals within a population may be in different states (e.g. different sizes or ages, living in different spatial locations) and these differences can have important consequences for population dynamics [18-24]. To account for how these differences influence persistence of interacting populations, we develop criteria for robust permanence for the dynamics of $k$ interacting structured populations. Roughly, these dynamics correspond to

replacing $x_i$ in (1) by vectors and $f_i(x)$ by matrices. In section 2, we describe these models in greater detail and review some basic concepts from dynamical systems theory. In section 3, we introduce the structured analog of average per-capita growth rates and define the concept of robustly unsaturated invariant sets. Necessary and sufficient conditions for an invariant set to be robustly unsaturated are proven. For Axiom A or uniquely ergodic invariant sets, the necessary and sufficient criteria are shown to be equivalent. In section 4, we use Morse decompositions and the criteria for robustly unsaturated invariant sets to develop necessary and sufficient conditions for robust permanence. In sections 5 and 6, we provide several applications of our results to spatially structured ecological models, structured epidemiological models, and models of gene networks. We apply the model to econometrics. To give an example investment needs inventory to be used towards the consummation of output, part of output is to be sold for consolidation of income, and some income might be replenished in the investment and so on. In finance a cash credit account could be debited or credited when there is say purchase of demand bills or discount of usance bills. Cash credit account is debited whenever cheques are issued towards the purchase of raw materials and so on. In fact some of the models have been applied to Organic psychiatry and we have shown how such models are useful in dealing with the propositions of Husserl, Derrida, Deleuze and others.

## 6. MODELS AND ASSUMPTIONS

Let $x_i$ denote the state of the $i$-th population that lies in the non-negative cone $C_i$ of $R^{ni.}$ Define $x=(x_1,...x_k)\epsilon C$ to be the non-negative cone of $R^n$ where $n=\sum_{i=1}^{k} n_i$. Let $C_i^+$ denote positive cone { $x_i \epsilon C_i$: $\prod_i$i noitalupop rof$^j$

$i$ *and* $C^+=\prod_i C^+_i$ , the positive cone for the interacting populations. If $x \mapsto A_i(x)$ is a map into $n_i x n_i$ matrices that describes the growth of population I, then the dynamics of the interacting populations are given by

$$(2)\frac{dx_i}{dt} = A_i(x)x_i \; i=1,...,k$$

Let $x.t$ denote the solution of (2), with the initial condition x, and more generally the semi flow generated by (2). Before stating our assumption on (2), recall a few definitions from dynamical systems. Given sets $I \subseteq R$ and $K \subseteq C$, let $K.I = \{x.t : t \epsilon I, x \epsilon K\}$. A set $K \subseteq C$ *is invariant* if K.t = K for all t>0, and *forward invariant* if $K.t \subset K$ for all $t > 0$. The *omega limit* set of a set $K \subseteq C$ equals $\omega(K) = \bigcap_{t \geq 0,} \overline{K.[t,\infty)}$. The *alpha limit* set of $K \subseteq C$ is $a(K) = \bigcap_{t \leq 0}$

$\overline{K.(-\infty,t]}$. Given a forward invariant set K, $B \subset K$ is an attractor for the semi flow *x.t restricted* to K provided there exists an open neighborhood $U \subseteq K$ of B such that $\omega(U) = B$. The stable set of a *compact invariant* set *K* is defined by

$$W^s(K) = \{x \in C : \omega(x) \neq \emptyset \text{ and } \omega(x) \subseteq K\}.$$

The semi flow generated by (2) is *dissipative* if there exists an attractor B with $W^s(B) = C$.

Throughout this paper, we make the following assumptions:

**A1**: $x \mapsto A_i(x)$ is continuous

**A2:** $A_i(x)$ is irreducible and has non-negative off- diagonal entries,

**A3:** *x.t* is defined for all $t \geq 0$

**A4:** *x.t* is dissipative with global attractor $\Gamma(A)$.

A1 is a basic regularity assumption. The non-negative of off- diagonal entries in A2 implies that there are no negative feedbacks between individuals of different states in population *i*. This assumption is meet for many types of structured models as discussed in sections 5 and 6. The irreducibility assumption A2 is generally meet and implies that all individuals within a population can pass through all states. In Remark1, we discuss how this irreducibility assumption can be relaxed. A3 ensures that the population dynamics are defined for all future time. A4 requires that population densities/ abundances eventually are uniformly bounded, a condition that should be meet for an biologically realistic model.

We say that the semi flow of (2) is *permanent* if this semi flow is dissipative and there is a positive attractor $B \subset C^+$ such that $W^s(B) \subset C^+$. We will denote this positive attractor by $\Lambda(A)$. Permanence ensures that populations can recover from rare large perturbations and allows for a diversity of dynamical behaviors.

For K $\subset$ C and $\delta > 0$, define the $\delta$- *neighborhood* of K as

$$N\delta(K) = \{x \in C : \backslash x - y < \delta \text{ for some } y \in K\}.$$

We define a $\delta$- *perturbation of (2)* to be a system of the form

$$\frac{dx}{dt} = \tilde{A}(x)x$$

that satisfies assumptions A1-A4, $\Gamma(\tilde{A}) \subset N$ *( $\Gamma(A)$), and* $||A(x) - \tilde{A}(x)|| \leq \delta$ *for all* $x \in N$ *($\Gamma(A)$).* (2) is *robustly permanent* if there exist >0 and $\epsilon > 0$ such that all $\delta$- perturbations of (2) are permanent, and $d(\mathcal{A}(\tilde{A}), C \backslash C^{+}) > \epsilon$ for all $\delta$- perturbations i.e., there is a common/ uniform region of repulsion around the boundary.

## INUNDATION OF COMPACT SETS

We begin by studying the linear skew product flows on $\Gamma(A) \times R^{n}_{i}$ defined by $(x,y).t = (x.t, B_i(t,x)y)$ where $Y(t) = B_i(t,x)$ is the solution to $\frac{dY}{dt}(t) = A_i(x.t)Y(t)$ with y(0) equal to the identify matrix. Our assumption that $A_i$ is irreducible with non- negative off diagonal entries implies that $B_i$ (t,x)$C_i \subset C_i^{+}$ for all $x$ and t > 0 (see, e.g., [25] ). A result of Ruelle [26, Prop, 3.2] implies that there exist continuous maps $u_i\, v_i : \Gamma(A) \rightarrow C_i^{+}$ with $u_i(x)| = |v_i(x)| = 1$ such that The line bundle $E_i(x)$ spanned by $u_i(x)$ is invariant i.e. $E_i(x.t) = B_i(x,t)E_i(x)$ for all t ≥ 0.The vector bundle $F_i(x)$ perpendicular to $v_i(x)$ is invariant i.e. $F_i(x.t) = B_i(x,t)F_i(x)$ for all $t \geq 0$.There exist constants $a > 0$ *and* $\beta > 0$ such that(3) $||B_i(t,x)F_i(x)|| \leq a\, exp(-\beta t)||B_i(t,x)E_i(x)||$ For all $x \in K$ *and* $t \geq 0$.**Interaction rates.** Given any x $\epsilon$ $\Gamma$(A), define the *interaction and inundation rate* of *species i at particle/proposition/problem/ state x as* $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right) in\ ||B_i(t,x)||$ when $x_i$ =0, $\lambda_i(x)$ provides an upper bound to the rate of growth of population *i* when introduced at infinitesimally small densities. Important properties of this invasion rate are summarized in the following two propositions. For instance, the second proposition implies if $(x.t)_i$ stays bounded away from zero (i.e. population *i* persist), then

$\lambda_i(x) = 0$. Alternatively, if $\lambda_i(x) < 0$,then population *i* is doomed to extinction. **Proposition 1.** $\lambda_i(x)$ *satisfies the following properties:* $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right) In\ |B_i\ (t,x)v|$ *for any* $v > 0$ *in* $R^{ni}$, *and* $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right) \int^{t}_{0} (A_i(x.s)u_i(x.s), u_i(x.s))ds.$

**Proof:** To prove the first property, we first show that

$$\lambda i(x) = \underset{t\to\infty}{limsup} \frac{1}{t} In\ |B_i(t,x)u_i(x)|.$$

To this end, let $v \in R^{ni}$ be any non-zero vector. Since $R^{ni} = E_{i(x)}$ $F_i(x)$, we can write $v = au_i(x) + w$ with $a \in R$ and $w \in F_i(x)$. Equation (3) implies

$|B_i(t,x)v| \leq a|B_i(t,x)u_i(x)| + |B_i(t,x)w| \leq |B_i(t,x)u_i(x)|(a+\alpha+exp(-\beta t)|w|)$. Hence, $\limsup_{t\to\infty}\frac{1}{t} In\, B_i(t,x)\, v| \leq$

$$\limsup_{t\to\infty}\frac{1}{t} In\, B_i(t,x)\, u_i(x) \leq \lambda_i(x)$$

for all non zero vectors $v \in R^{ni}$. Since $||B_i(t,x)|| = sup|v|=1 B_i(t,x)v|$, this inequality implies that $\lambda i(x) = \limsup_{t\to\infty}\frac{1}{t}$ In $|B_i(t,x)v|$, this inequality implies that $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right)$ In $|t,x)u_i(x)|$. Now let v $\epsilon$ $C_i^+$. Then, we can write $v = au_i(x) + w$ with $a > 0$ and $w \in F_i$.

Equation (3) implies

$$|B_i(t,x)v| \geq a|B_i(t,x)u_i(x)| - |B_i(t,x)w|$$

$$\geq |B_i(t,x)u_i(x)|(a-\alpha+exp(-\beta t)|w|).$$

Since a>0, this inequality implies that

$$\lambda i(x) \geq \limsup_{t\to\infty}\frac{1}{t} \text{In } |B_i(t,x)v| \geq \lambda i(x) \geq \limsup_{t\to\infty}\frac{1}{t} \text{In } |B_i(t,x)u_i(x)| = \lambda i(x)$$

To prove the second assertion, let $b_i(t,x) = In\, |B_i(t,x)u_i(x)|$. Invariance of $u_i(x)$ implies that $b_i(t,x)$ is additive:

$$B_i(t+s,x) = In|B_i(t+s,x)\, u_i(x)| = In\, |B_i(s,x.t)B_i(t,x)u_i(x)$$

$$= \text{In } |B_i(s,x.t)\, u_i(x.t)||\, B_i(t,x)\, u_i(x)|$$

$$= \text{In } b_i(s,t.x)\, b_i(t,x)$$

Additivity of $b_i(t,x)$ and the fact that $b_i(0,x) = 0$ implies

$$\frac{d}{dt}b_i(t,x) = \lim_{s\to 0}\frac{bi(t+s,x-bi(t,x)}{S-}$$

$$= \lim_{s\to 0}\frac{bi(s,x.t)}{S} = \frac{d}{ds}|_{s=0}\, b_i(s,x.t)$$

$$=$$

$$=(A_i(x.t)u_i(x.s),u_i(x.s),$$

The Fundamental Theorem of Calculus implies

$$B_i(t,x)= \int_0^t (A_i(x.s)\, u_i(x.s), u_i(x.s))ds\backslash$$

and the second assertion follows.

**Proposition 2:** *For the solutions of (2) we have:* if $x_i>0$, then $\lambda i(x) \leq 0$. If $\lambda i(x)<0$, then $\lim_{t\to\infty} (x.t)_i=0$ If $limsup_{t\to\infty}$ $(x.t)_i|>0$, then $\lambda i(x)=0$.

**Proof:** First, assume that x satisfies $x_i>0$. Since the semi flow is dissipative, there exists $\gamma >0$ such that $|x.t| \leq \gamma$ for all $t \geq 0$. Proposition 1 and the definition of the skew product flow imply

$$\lambda i(x) = \limsup_{t\to\infty} \frac{1}{t} \text{In } |B_i(t,x)x_i| = \limsup_{t\to\infty} \frac{1}{t} \text{In } |(x.t)_i|$$

$$\leq \limsup_{t\to\infty} \frac{In\,\gamma}{t} = 0$$

To prove the second assertion, assume that $\lambda i(x) < 0$. $If\ x_i=0$, then the invariance of the faces of C imply that $(x.t)_i=0$ for all $t\geq 0$. Alternatively, if $x_i> 0$, then Proposition I and the skew product flow imply

$$\limsup_{t\to\infty} \frac{1}{t} In|(x.t)_i| = \lambda i(x) < 0.$$

The final assertion follows from the first two assertions. 5

**3.2. Invariant measures.** We review some definitions from ergodic theory. Given a Borel probability measure $\mu$ on ***C***, the *support* of $\mu$, denoted supp( $\mu$),is the smallest closed set whose complement has measure zero. A Borel probability measure $\mu$ is called *invariant* for (2) provided that Integral $h(x)d\mu(x) =$ Integral $h(x.t)d\,\mu(x)$for all $t \geq 0$ and for all bounded continuous functions $h$:C→R. An invariant measure $\mu$ is called *ergodic* provided that $\mu(B)$=0 or 1 for any invariant Borel set *B.*

For an invariant measure $\mu$ define the invasion rate species $i$ with respect to $\mu$ as

$$\lambda_i(\mu) = \int_c \langle A_i(x)u_i(x), u_i(x)\rangle d\mu(x)$$

**Proposition 3:** Let $\mu$ be an *invariant measure.* $\lambda_i(\mu)$ *Satisfies the following properties:* $\lambda_i(x) = \lim_{t\to\infty} \frac{1}{t}\int_0^t \langle A_i(x.s)u_i(x.s), u_i(x.s)\rangle, ds$ exists for $\mu$ *- almost every* $x$. *Moreover, if* $\mu$ *is ergodic,* $\lambda_i(x) = \lambda_i(\mu)$ for $\mu$-*almost every* $x$.*If* $\mu$ *is ergodic, then there exists* $I \subset \{1, \dots, k\}$such that $\mu(\Pi_{i\in I} C_i^+)$=1 and $\lambda_i(\mu) = 0$ *for all* $i \in I$.

**Proof:** Let $\mu$ be an invariant measure. Define $h_i(x) = \langle A_i(x)u_i(x), u_i(x)\rangle$. The Birkhoff's' ergodic theorem and Proposition 1 imply

$$\lambda_i(x) = \lim_{t\to\infty} \frac{1}{t}\int_0^t h_i(x.s)ds$$

Exists for $\mu$- almost every $x$. Moreover, $\lambda_i(x) = \lambda_i(\mu)$ $\mu$- almost surely if $\mu$ is ergodic.

Assume $\mu$ is ergodic. By ergodicity and invariance of the faces of **C**, there exists $I \subset \{1, \dots, k\}$ such that $\mu(\Pi_{i\in I} C_i^+)$=1. Let $K \subset \Pi_{i\in I} C_i^+$ be a compact set such that $\mu(K) > 0$. The Poincare recurrence theorem and the Birkhoff ergodic theorem imply that there is $x \in \Pi_{i\in I} C_i^+$such that $x.t_n \in K$ for some $t_n \uparrow \infty$ and $\lambda_i(x) = \lambda_i(\mu)$ for all $i$. The third assertion of Proposition 2 implies $\lambda_i(\mu) = 0$ for all $i \in I$.

**3.3. Robustly unsaturated sets**. Let $K \subset \partial C$ be a compact isolated invariant set for the flow $x.t$ restricted to $\partial C$.*K* is unsaturated if $W^s(K) \subset \partial C$ and K is isolated for $x.t$. If *K* is not unsaturated, then *K* is saturated. *K* is robustly unsaturated for (2) if there exists $\delta > 0$ such that the continuation of K for any $\delta -$ perturbation of (2) is unsaturated.**Theorem 1:** Let *K* be a compact isolated invariant set for $x.t$ restricted to$\partial C$. If one of the following equivalent conditions hold For all invariant measures $\mu$ supported by *K* (4) $\max_{1\le i\le k} \lambda_i(\mu) = 0$

There exist $p_i > 0$ such that (5) $\sum_{1\le i\le k} p_i \lambda_i(\mu) = 0$

For all ergodic probability measures supported by *K*.

then *K* is robustly unsaturated for (2). Alternatively, if $x \to A(x)$ is twice continuously differentiable and *K* is robustly unsaturated, then (4) holds for all ergodic measures $\mu$ supported by *K*.

**Remark 1**: For some applications (e.g., the disease model considered in section 6.1) it useful to relax the irreducibility assumption A2. For instance, if there exists an open neighborhood U of K such that for each $i$, $A_i(x)$ has a fixed off-diagonal sign pattern for all $x \in U$,then $A_i(x)$ can be decomposed into a finite number, say $m_i$, of

irreducible components. For each of these irreducible components, one can define $\lambda_i^j(\mu) = \int_C \langle A_i^j(x)u_i^j(x), u_i^j(x)\rangle \, d\mu(x)$ where $A_i^j(x)$ is the submatrix of $A_i(x)$ corresponding to the $j$-th irreducible component of $A_i(x)$ *and* $u_i^j(x)$ is the continuous invariant sub bundle for the irreducible linear cocycle determined by $A_i^j(x)$. If we define $\lambda_i(\mu) = \max_{1\le j\le m_i} \lambda_i^j(\mu)$, then all of the assertions of Theorem 1 still hold.

**Proof:** To prove the first assertion of the theorem, we need the following lemma. We call an invariant measure $\mu$ for (2) saturated if $\lambda_i(\mu) \le 0$ for all $i$

**Lemma 1**: Let $\delta_n$ be a non –negative sequence that converges to zero as $n\to\infty$. If $\mu_n$ are saturated invariant measures for $\delta_n -$ perturbations of (2), then the weak$^*$ limit points of $\{\mu_n\}_{n=1}^{\infty}$ is a non-empty set consisting of saturated invariant measures for (2).

**Proof:** Let $\delta_n$ be a non-negative sequence that coverages to zero. Let $\mu_n$ be saturated invariant measures for $\delta_n$ perturbations, $\dot{x} = A^n(x)x$, of $\dot{x} = A(x)x$. By weak$^*$ compactness of Borel probability measures supported on $N_1(\Gamma(A))$, there exist weak$^*$ limit points of $\mu_n$ . Let $\mu$ be such a weak$^*$ limit point. Since $\mu_n$ are supported by $N_1(\Gamma(A))$, for all $n$ sufficiently large, $\mu$ is supported by $N_1(\Gamma(A))$.To verify that $\mu$ is an invariant measure for (2),let $h: C \to R$ be a bounded continuous function. Let $t > 0$ and $\in > 0$ be given. Choose $n$ sufficiently large so that $|\int_c h(x)\, d\mu(x) - \int_c h(x) d\mu_n(x)| \le \in$, $|\int_c h(x.t)\, d\mu(x) - \int_c h(x.t) d\mu_n(x)| \le \in$, and $|h(x.t) - h(x^n.t)| \le \in$ for all $x \in N_1(\Gamma(A))$.Then $|\int_c h(x.t) - h(x) d\mu(x)|$ is

$\le \quad |\int_c h(x.t)\, d\mu(x) - \int_c h(x.t)\, d\mu_n(x)| + |\int_c h(x.t)\, d\mu_n(x) - \int_c h(x)\, d\mu(x)|$

$\le \quad \in + \int_c |h(x.t) - h(x^n.t)|\, d\mu_n(x)| + |\int_c h(x^n.t)|\, d\mu_n(x) - \int_c h(x) d\mu(x)|$

$\le \quad 2 \in + |\int_c h(x^n.t) - h(x)\, d\mu_n(x)| + |\int_c h(x) d\mu_n(x) - \int_c h(x) d\mu(x)|$

$\le \quad 3 \in$

Where the last line follows from the invariance of $\mu_n$ for $x^n.t$. Since $\in > 0$ and t > 0 are arbitrary, we have $\int_c h(x.t)d\mu(x) = \int_c h(x)d\mu(x)$ for all t >0 and all bounded continuous functions $h: C \to R$. It follows that $\mu$ is an invariant measure for (2). To see that $\mu$ is saturated, define $h_i^n(x) = \langle A_i^n(x)u_i^n(x), u_i^n(x)\rangle$ where $u_i^n(x)$ spans the invariant one-dimensional bundle given by $[26, Prop.3.2]$ for $\frac{dx}{dt} = A^n(x)x$. Since $A_i^n(x) \to A_i(x)$ uniformly for $x \in N_1((A))$ as $n \to \infty$, $[26, Prop.3.2]$ implies that $u_i^n(x) \to u_i(x)$ converges uniformly for $x \in N_1(\Lambda(A))$. Given $\in > 0$, choose $n$ sufficiently large so that $|h_i^n(x) - h_i(x)| \leq \in$ for all $x \in N_1(\Lambda(A))$ and $|\int_c h_i^n(x)\,d\mu(x) - \int_c h_i^n(x)\,d\mu_n(x)| \leq \in$.

Then

$$\lambda_i(\mu) \leq \int_c |h_i(x) - h_i^n(x)|d\mu(x) + |\int_c h_i^n(x)\,d\mu(x) - \int_c h_i^n(x)\,d\mu_n(x)|$$

$$+ \int_c h_i^n(x)\,d\mu_n(x)$$

$$\leq 2\in$$

Since $\in > 0$ is arbitrary, $\lambda_i(\mu) \leq 0$ for all $i$ and $\mu$ is saturated.

Using this lemma, we prove that if K is saturated, then there exists a saturated invariant measure $\mu$ supported by K. Assume K is saturated. Work of Hofbaue and So $[7, Thm.2.1]$ implies that either there exists $y \in W^s(K) \cap C^+$ or K is not isolated for the unrestricted flow $x.t$ If there exists $y \in W^s(K) \cap C^+$, then for all $t \leq 0$ define $v_t = \frac{1}{t}\int_0^t \delta_{y.s}\,ds$ where $\delta_{y.s}$ denotes a Dirac measure based at the point $y.s$. Dissipativeness of (2) and weak* compactness of the Borel probability measures supported by (A) imply there exist $t_k \to \infty$ such that the sequence $v_{tk}$ converges in the weak* topology to a Borel probability measure $\mu$ with support in K.A standard argument implies that $\mu$ is $x.t$ invariant. Define $h_i: C \to R$ by $h_i(x) = (A_i(x)u_i(x), u_i(x))$.Proposition 2 and weak* convergence imply that

$$0 \geq (\lambda i(y) = \limsup_{t\to\infty} \frac{I}{t}\int_0^t h_i(y.s)ds$$

$$= \limsup_{t\to\infty} \int_c h_i\,dv_t$$

$$\geq \limsup_{t\to\infty} \int_p h_i\,d\,v_{tk} = \lambda i(\mu)$$

For all $i$. Alternatively suppose that K is not isolated for the semi flow Then there exists a sequence of positive $\omega - limit$ sets that accumulate on K. Let $\mu_n$ be a sequence of ergodic probability measures supported by these $\omega -$ limit sets. Propositions 1 and 3 imply that $\mu_n$ are saturated for all $n$ Applying Lemma I with $\delta_n$=0 for all $n$ implies that there exists a saturated invariant measure $\mu$ suppoorted by *K*. Hence we have shown that if *K* is saturated then there exists a saturated invariant probability measure $\mu$ supported by *K* Equivalently, (4) holding for all invariant probability measures with support in *K* implies that *K* is unsaturated.

Next we show that if *K* not robustly saturated, then there exists a saturated invariant measure supported by *K* Indeed suppose *K* is not robustly saturated, Then there exists a non-negative sequence $\delta_n$ converging to zero and a sequence of saturated measures $\mu_n$ for $\delta_n$-perturbations of (2)with support in the continuation of *K*. Let $\mu$ be a weak$^*$ limit point of $\{\mu_n\}\infty_{n=1}$ Lemma 1 implies that $\mu$ is saturated. Moreover, since the continuation of *K* converges to *K* as $\delta_n \to 0, \mu$ is a saturated invariant measure for (2) supported by *K* Hence we have shown that if *K* is not robustly saturated then there exists a saturated invariant measure supported by *K* Equivalently if (4) holds for all invariant measures with support in *K* then *K* is robustly unsaturated.

To see the equivalence of the conditions given by (4) and (5) let $\Delta = \{p \in R_+^k : \sum_i P_i = 1\}$ and notice that

$$\text{Min max } \lambda i\ (\mu = \min_{\mu} \max_{p \epsilon \Delta} \sum_i Pi\lambda i(\mu)$$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K* The Minimax theorem (see, e. g.,[27] implies that

$$(6)\ \min \max \lambda i(\mu) = \min_{\mu} \max_{p \epsilon \Delta} \sum_i Pi\lambda i(\mu)$$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K*. Since min$\mu$

$\sum_i P_i \lambda_i(\mu)$ is attained at an ergodic probability measure with support in *K*, the equivalence of the conditions given by (4) and (5) is established.

To prove the final assertion of the Theorem assume $x \to A(x)$ is twice continuously differentiable and there exists a saturated ergodic measure $\mu$ supported by *K* Proposition 3 implies that there exists *I*C {1,…,k} such that (□$i \in iC_i^+$)=1and $\lambda i(\mu) = 0$ For all $i \in I. SinceKc\,.\partial$□,$\{1, \ldots . k\}\backslash I$ is non-empty. We will show that (2) is not robustly permanent by proving that for all $\delta > 0$ there exists a $\delta$ −perturbation of (2) for which the continuation of *K* is saturated Let $\delta > 0$ be given Choose V $c$ N$\delta$ (*Γ*(A)) to be a compact neighborhood of Γ(A) such that V.t $c$ int *V* for all t > 0.Let $\eta > 0$ be such that *W*=:*N* $\eta$(Γ(A))$c$ *V* Let p:c→ $[0,1] be$ a$C^{\infty}$function such that $p(x) = 1 \quad for\ all\ x \in \Gamma(A) and\ p(x) = 0 for\ all\ x \in C\backslash W\ Define\ \tilde{A} = (\tilde{A}_1, \ldots, \tilde{A}_k)$ by

$$\tilde{A}_i(x) = \{ \frac{A(x)}{A(x) - \frac{\delta}{2Idp\,(x)}} \quad \frac{if\ i \epsilon I}{if\ i \notin I}$$

Where Id denotes the identity matrix of appropriate dimension Let x. t denote the semi flow of $x = \tilde{A}(x)x$.Since$x.t = x.t$ $whenever\ x.[0,t] \in C\backslash W, it\ follows\ that \quad \tilde{V}.t \subset intV\ for\ all\ t > 0. Hence, \Gamma(\tilde{A})c\ V\ c\ N\delta\ (\ (A))). We\ also\ have\ ||A(x) - \tilde{A}(x)|| \leq \frac{\delta}{2}$Therefore $x = \tilde{A}(x) is\ a\ \delta -$ perturbation of (2) By construcation,$x.t = x.t\ for\ all\ x \in K and\ t \geq 0. Hence, \mu$ is ergodic for the semiflow of $\dot{x} = \tilde{A}(x)x\ and\ \lambda_i(\mu) \leq -\frac{\delta}{2}$ for this semiflow and $i \in \{1, \ldots, k\}\backslash I$.Let $\mathcal{L}$ and O be the Lyapunov exponents and Oseledec regular points supported by $\mu$ for $x = \tilde{A}(x)x(see, e.g.,$

[28] for definitions).At each point $x \in O$, the splitting of $R^n$ determines three subspaces: the stable subspace $E^s(x)$, the center subspace $E^c(x)$, and the unstable subspace $E^u(x)$ Proposition 3 and our choice of Ã imply that $E^s(x) \cap C^+ \neq \emptyset$. The Pesin stable manifold theorem [28, Corollaries 3.17 and 3.18] implies that tangent to $E^s(x)$ $E^c(x)$ and $E^u(x)$ are locally $x.t$ −invariant families of $C^1$ discs $W^s_x, W^c_x$ and $W^u_x$ corresponding to the stable ,center and unstable manifolds The family of stable manifolds $W^s_x$ is contained in $W^s$ (Ã,supp $(\mu)$). $Since\ E^s(x) \cap C^+ \neq \emptyset$, $W^s_x \cap C^+ \neq \emptyset$ for some $x \in K$. Consequently *K* is saturated for $x = \tilde{A}(x)x$.

The difference between the sufficient and necessary condition is that for the sufficient condition (4) has to hold for all invariant measures supported by *K*, while for the necessary condition (4) has to hold only for ergodic measures supported by *K*. since the invariant measures lie in the convex hull of the ergodic measures the sufficient condition can be more restrictive then the necessary condition.

For three Classes of invariant sets the necessary and sufficient conditions coincide Recall a compact invariant set *K* for (2) is uniquely ergodic if K only supports one invariant measure Recall a compact invariant set *K* is Axiom A if the flow of (2) restricted to *K* is transitive and hyperbolic (see, e.g.,[29] for definitions.)

**Corollary 1:** Assume $x \to A(x)$ is twice continuously differentiable It K$\subset \partial$□ is a compact invariant set and either is uniquely ergodic or Axiom A, then *K* is robustly unsaturated if and only if (4) holds for all ergodic, measures $\mu$ with support in *K*

Proof. If K is uniquely ergodic then the assertion follows immediately For an Axiom A invariant set Sigmund [30, thm.1] has proven that invariant measures supported by periodic orbits of *K* are dense in the set of invariant probability measures supported by *K* In particular any invariant measure can be approximated by an ergodic measure and consequently the result follows.

Another special case where the necessary and sufficient conditions coincide is when *K* supports all populations except one for discrete time models this case was considered by [22].

**Corollary2:** Assume $x \to A(x)$is twice continuously differentiable Let K$c$ $II_i \geq 2C_i^+$ be a compact invariant set. Then K is robustly unsaturated if and only if $\lambda_1(\mu) > 0$ for all invariant measures supported by *K*

## 4. MORES DECOMPOSITIONS AND ROBUST PERMANENCE

To state the sufficient condition for robust permanence, we use a characterization of permanence due to Garay [8] and Hofbauer and So [7] that involves Morse decompositions of the boundary flow Conley[31] defined a collection of sets$\{M_{1,\ldots,}M_m\}$ to be a Morse decomposition for a compact invariant set *K* if $M_1\ldots,M_m$ are pairwise disjoint, compact isolated invariant sent for the flow of (2) restricted to K. For each $x \in K$ there are integers r=r$(x) \leq s =$

$s(x)$ such that $a(x) \subseteq M_r$and $\omega(x) \subseteq M_s$. If $r(x) = s(x), then\ x \in M_r$. Garay, Hofbauer and so [7,8] proved the following characterization of permamnce

**Theorem 2** (Garay, Hofbauer-so) if $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$ then (2) is permanent if and only if each of the components $M_1$ are unsaturated.

Theorem 1 and 2 imply the following result: As stated in the above all these results are true for most of problems, interactions, transactions in cosmology, econometrics, finance and other subjects, which we shall illustrate as we go on constructing models for the same.

**Theorem 3:** If $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$ and (4) holds for each of the components of the More decomposition, then (2) is robustly permanent Conversely if $x \to A(x) is\ twice\ continuously\ differentiable$

*and (4) is violated by an ergodic measure supported by one of the components of the Morse decomposition, then (2) is not robustly permanent.* Theorem 3 in conjunction with Corollaries 1 and 2 yield a characterization of robust permanence for a class of structured models.

**Corollary 3:** *Let $\{M_1,\ldots..,M_m\}$ os a Morse decomposition for* $\partial$C. *Assume $x \mapsto A(x)$ is twise continuously differentiable and for each Morse component $M_i$ one of the following assertions hold $M_i$ is Axiom A, $M_i$ is uniquely ergodic, or There exists $j \in \{1,\ldots,k\}$ such that $M_i$ ----- $\{x \in C : x_l \gg 0 \text{ for } l \neq j\}$ Then (2) is robustly permanent if and only if (4) holds for all ergodic measures supported by $U_i M_i$.*

Our result also provided a structured analogue for characterizing totally permanent systems[32]. For an ergodic probability measure$\mu$, define sup ($\mu$) to be the subset $I$ -----$\{1,\ldots,k\}$ such that $\mu(---_{i \in I} C_+) = 1$

**Corollary 4:** *The following statements are equivalent: For all ergodic probability measures $\mu$ with support in* $\partial$C, *$\lambda_i(\mu) > 0$ for all $I \in \{1,\ldots,k\} \setminus supp(\mu)$ (2) and all of its subsystems are robustly permanent.*

*Proof. Suppose the first statement holds. Then unsaturated condition given by (5) with p= (1,…,1) and Theorem 3 implies that (2) and all its subsystems are robustly permanent.. The other direction follows form the necessary condition for robust permanence in Theorem 3 and the second assertion of Proposition 3.*

*To illustrate the broad applicability of theorem 3, we develop applications to spatially structured models in section 5 and to a disease model and a gene network model in section 6.*

## 5. APPLICATIONS: PATCH MODELS

A fundamental application of our result are to spatially structured models, with *k* species dispersing between *m* patches {33-35}:

(7) $\frac{dx_i^j}{dt} = x_i^j f_i^j(x^j) + \sum d_i^{jl} x_i^l - e_i^j x_i^j$ i = 1,..,k j=1,…,m

*l*

Where $x_i^j$ denotes the density of species I in patch *j*, $x^j = (x^j_1,.....,x^j_k)$ is the vector of species densities in patch *j*, $f^j_i$ is the per-capita growth rate of species *i* in patch *j*, $d_i^{jl} \geq 0$ is the dispersal rate for species *i* from patch *l* into patch *j*, and $e_i^j$ is the emigration rate of species *i* out of patch *j*. Hence $e_i^j \geq \sum_l d_i^{lj}$. We assume that the matrices $(d_i^{lj})_{j,l}$ are irreducible for each species *i* . it is then easy to write (7) in the form (2), with $x_i = (x^l_i, ..., x_i^m)$ 'where' denotes transpose.

For a single species (7) generates a monotone flow. Under mild assumption, e.g., each $f^j_i$ *is* decreasing and negative for large densities, there is a globally stable equilibrium, see, e.g.,[36] and [37, sect. 5.4]. This equilibrium will be positive if the invasion rate at the origin λ (0), which is given by the leading eigenvalue (stability modulus) of the matrix $(d^{jl} + f^j(0) - e^j)$, is positive.

5.1 Two types of debitable and creditable systems:. For *two competing species*, e.g., $f^j_i$ is decreasing with respect to both species and is negative at large densities, the dynamics are still monotone, and hence almost all orbits converge to an equilibrium. Robust permanence requires two conditions. First, $\lambda_i$ (0) >0 for both species in which case there are two single species equilibria $E_1$ and $E_2$. Second, the invasion rates $\lambda_2(E_1)$ and $\lambda_2(E_2)$ are positive. Whether there is a unique positive globally stable equilibrium for the robustly permanent system depends in a delicate way on the

system parameters, see [38]. For slowly dispersing populations (i.e. $0< e^j_i \ll 1$ for all *i, j* ), the robust permanence condition is particularly

Straight forward to verify. Let $E^j_i$ be the largest solution to $x^j_i \quad f_i{}^j \ (x^j_i) = 0$ i.e. the equilibrium attained by species *I* in patch *j* when there is no dispersal and no competitors. For slowly dispersing populations, $E_1$ is close to the product $(E_1{}^1, \ldots, E_1{}^m)$, and the matrix $A^2(E_1)$ is close to the diagonal matrix $\text{diag}(f^i{}_2(E_1{}^j))^m_{j=1.}$ Hence the invasion rate $\lambda_2(E_1)$ is close to $\max_j f_2^j(E^j_1)>0$. So we obtain: The system with two slowly dispersing competing species is robustly permanent if and only if system refers to system of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

(8) $\max_j f_2^j(E^j_1)>0$ and $\max_j f_2^j(E^j_2)>0$

Intuitively, robust permanence requires for each species there is at least one patch where it can persist.

For *systems, particle interactions, financial transactions, cosmological interactions and reactions, as we shall show model by model the* situation are even easier. Since (under mild assumption on the prey dynamics) the Morse decomposition of $\partial C$ is given by two equilibria $\{0,E_1\}$, the spatial predator-prey system is robustly permanent, if the predator can invade at the prey equilibrium: $\lambda_2(E_1) > 0$. However, the global dynamic are likely to be more complicated. For instance, if in each patch there is a globally stable limit circle (e.g.,[39-41]) and dispersal is sufficiently slow (i.e. $0< e^j{}_i \ll 1$for all *i, j*), our result about robust permanence and the theory of normally hyperbolic manifolds [42] imply that there is a positive *m* dimensional torus, which attracts almost all orbits in C, whenever $\max_j f_2^j(E^j_1)>0$

**5.2 Rock-Paper-Scissors Dynamics.** The Lotka-Volterra model of rock-paper-scissor dynamics is a simple model that is used as property for understanding intransitive ecological outcomes [43,44]. Here, a simple spatial version of this dynamic is given by

$$\frac{dx_1^j}{dt} = x_1^j(1 - x_1^j - \beta^j \ x_2^j \ - \alpha^j \ x_3^j) + \sum d_1^{jk} \ x_1^k \ - e_1^j \ x_1^j$$

$\frac{dx_2^j}{dt} = x_2^j(1-\alpha^j x_1^j - x_2^j - \beta^j x_3^j) + \sum d_2^{jk} x_2^k - e_2^j x_2^j$

$\frac{dx_3^j}{dt} = x_i^j(1-\beta^j x_1^j - \alpha^j x_2^j - \beta^j x_3^j) + \sum d_3^{jk} x_3^k - e_3^j x_3^j$

Where $a^j \in (0,1)$ and $\beta^j > 1$ for all $j$. A more general version of this dynamic is presented in [44]. Under the assumption that there is no cost to dispersal (i.e. $e^j_i = \sum_l d^{l,j}_i$ for all $i, j$) , the maximal invariant set in $\partial C$ consist of the origin 0 and a heteroclinic cycle connecting positive single species equilibria $E_1$, $E_2$ and $E_3$. For slowly dispersing population ($0 < e^j_i \ll 1$ for all $i$, j), $E_1$ is close to $(x_1, x_2, x_3) = (1,\ldots,1,0,\ldots,0,0,\ldots,0)$ and the invasion rates $\lambda_2(E_1)$and $\lambda_3(E_1)$are close to $\max_j 1-a^j$ and $\max_j 1-\beta^j$. For the other equilibria $E_i$, the debitable rates of the missing species are also given by $\max_j 1-a^j$ and $\max_j 1-\beta^j$. Consider the Morse decomposition of $\partial C$ given by 0 and the heteroclinic cycle. Since $\lambda i(0) = 1 > 0$ for $i = 1, 2, 3, 0$ is robustly insaturated. An algebraic computation reveals that (4) holds for all invariant supported by the heteroclinic cycle (i.e. all convex combination of the Dirac measures supported by the equilibria) if and only if

$$\max_j(1-a^j) > -\max_j(1-\beta^j)$$

Equivalently,

$$2 > \min_j a^j + \min_j \beta^j$$

In particular, even if the heteroclinic cycle is attracting for each patch when the system is uncoupled (i.e. $2 < a^j + \beta^j$ for all $j$ ), it can be repelling for the weakly coupled system.

5.3. Three species **Lotka Volterra in spatially homogenous environments.** Consider Lotka-Volterra dynamics in a spatially homogenous environment: $(x+a)^n =$

(9) $$\frac{dx_i^j}{dt} = x_i^j\left(r_i - \sum_s a_{is}x_s^j\right) + \sum_l d_i^{jl}x_i^l - e_i^j x_i^j, \qquad i = 1,\ldots,k;\ j = 1,\ldots,m$$

We assume that dissipative can be shown by a linear Liapunov function: There are $c_i > 0$ such that $\sum_i c_i\, a_{is} x_i x_s < 0$ holds for all $x \in R_+^k \setminus \{0\}$. Then the weighted sum of densities across all patches $\sum_{i,j} c_i\, x_i^j$ is decreasing for large densitities and hence (9) generates a dissipative semiflow with global attractor $r$. We further assume that for each patch *j* and each species *i* the immigration rate equals the emigration rate , i. e,

(10) $$e_i^j = \sum_l d_i^{jl}$$

This guarantees that the set $H = \{x \in C: x_i^j = x_i^l \forall i, j, l\}$ of spatially homogenous states is forward invariant under (9).

Now consider the Lotka-Volterra dynamics without spatial structure

(11) $$\frac{dx_i}{dt} = x_i(r_i - \sum_s a_{is} x_s), \qquad i = 1, \ldots, k$$

(11) $$\frac{dx_i}{dt} = x_i(r_i - a_{is} x_s), \quad i = 1, \ldots, k$$

It has been shown [12,ch.16.1,16.2] That for *k* =3,(11) is robustly permanent (robust meaning here within the class of Lotka –Volterra systems) if and only if all equilibria on the boundary of $R_+^3$ are unsaturated and whenever there is a heteroclinic cycle connecting the one species equilibria then this cycle is repelling. Equivalently the following four conditions hold. There exists an interior equilibrium $\dot{x}$ (i.e $A\dot{x} = r \;\; with\; \dot{x} \gg 0$); Det (-A)>0; the 2 species subsystems are not bistable competition systems if there is a heteroclinic cycle between the one species equilibria, say $\mathbf{E_1 \to E_2 \to E_3 \to E_1}$ then the following inequality holds

12) $$\lambda_2(E_1)\lambda_3(E_2)\lambda_1(E_3) > |\lambda_3(E_1)\lambda_1(E_2)\lambda_2(E_3)|$$

Furthermore, the proof in [12, ch.16.1, 16.2] shows that the boundary flow of (11) has a simple Morse decomposition: if (iv) applies then there are two Morse sets, the origin 0 is a repeller, and the heteroclinic cycle is

the dual attractor (within$\partial R^3_+$); if (iv) does not apply then the finitely many boundary equilibria form a Morse decomposition of the boundary flow.

**Species, population we mean types of and population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.**

**Theorem 4:** Under the above assumptions,(9) for k=3 is permanent if and only if *(i),(ii),(iii)* and *(iv)* hold.

Proof: Suppose (9) is robustly permanent. Since the nonspatial system (11) is the restriction of (9) to the invariant subspace **H**, (11) must be robustly permanent and hence*(i)-(iv)* hold

Conversely, suppose that (i)-(iv) hold. We show that the above Morse decomposition of the boundary flow of (11) (now in **H**) is also Morse decomposition for (9) restricted to the maximal invariant subset of $\partial C$ which is contained in the spatial two species subsystems of (9). For this we show that in each 2 species system the orbits of (9) converge to **H** (For the one species subsystems this is obvious.) If a two species subsystem of (11) has an internal equilibrium, say $\mathbf{E_{12}}$ then by assumption (ii),this is not a saddle but a sink, and hence $a_{11}a_{22} > a_{12}a_{21}$. This implies that the 1-2 submatrix of A is VL-stable [12,ch.15.3,],and [45] implies that $\mathbf{E_{12}}$ is the global attractor for the spatially structured system (9) restricted to the first two species.

If a two species subsystem of (11) has no internal equilibrium, then a one species equilibrium say $\mathbf{E_1}$ is the global attractor. Then Lemma 2 (see below) show the corresponding result for (9).A similar proof applies to the case where the origin **0** is the global attractor of a two species subsystem.

Finally observe that under the above assumptions (9)-(10), the invasion rates $\lambda_i$ at spatially homogenous boundary equilibria are the same for (9) and (11) Since all boundary equilibria are unsaturated for (11), they are for (9) as well. The same applies in case (iv) to the heteroclinic cycle. Note that the condition (12) on the eigenvalues is equivalent to the fact that all invariant measures supported on the heteroclinic cycle are unsaturated, see [16,Ex.4.5]

Species we mean types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

**Lemma 2:** Suppose, in a two species Lotka –Volterra system (11) (k=2) all interior orbits converge to the one species equilibrium $\mathbf{E_1}$ .Then all orbits of the spatial version (9) (k=2) converge to the spatially homogeneous one species equibrium $\mathbf{E_1} \in \mathbf{H}$.

**Proof:** If the system is of predator-prey or mutualistic type then Hasting's [45] result applies again. So let us assume that the local interaction is given by the competition system

(13) $\frac{dx_1}{dt} = r_1 x_1(1 - x_1 - ax_2)$ $\frac{dx_2}{dt} = r_2 x_2(1 - \beta x_1 - x_2)$

With 0<a <1< $beta$.In this case, Kazuo Kishimoto (letter to JH, Oct 1987) has given the following argument. There exists a family of forward invariant rectangles contracting to $\mathbf{E_1}$. This shows that all interior solutions of (9) (k=2) converge to $\mathbf{E_1}$.

Theorem 4 shows that a spatial network of identical patches is permanent, if the within patch dynamics is permanent, and three species Lotka-Volterra dynamics. It is not clear how to extend this result to more than three species. The crucial step in the three species case is that permanence precludes bistable two species subsystems. The spatial version of a bistable two species system allows plenty of stable equilibria outside **H,** see [33].However, permanent four species Lotka Volterra systems may have two species bistable subsystems [12,ch.16.4].So there may be invariant sets on the boundary outside **H** which need to be unsaturated to make (9) permanent.

## 6. MORE APPLICATIONS

6.1. **Disease, particle, product, marketing, financial, econometric dynamics with density-dependent demography**. GAO and Hethcote [46] introduced a model of disease dynamics with density-dependent demography. Here we consider a variation of their model in which a population of size P has a per-capita birth rate b(P) and per –capita death rate d. To ensure that the population persists in the absence of the disease, we assume that there exists K>0 such that b(K)=d. To describe the disease dynamics, let *S* be the number of individuals susceptible to the disease, *I* the Number of infected individuals, and R the number of individuals that have recovered from the disease. We assume that P=S+*I* +R. On the same lines we can think of sick and non performance accounts, and those which are conducting themselves well, and those which are satisfactory and find a parallel formula. The same thing holds true in case of a black hole eating matter and black hole radiating radiation.. If the disease transmission is asymptotic, β is the contact rate between susceptible and infected individuals, γ is the rate at which individuals recover from the disease, and m the mortality rate due to disease, then the population dynamics is given by

$$\frac{dp}{dt} = b(P)P - dP - mI$$

$$\frac{dI}{dt} = \beta(P - I - R)\frac{I}{\epsilon + P} - (\gamma + d + m)I$$

$$\frac{dR}{dt} = \gamma\, I - dR$$

Where $\epsilon>0$ is a constant It is useful to introduce a change of coordinates in which $y=\frac{I}{P}$ and $z=\frac{R}{P}$:

$$\frac{dP}{dt} = (b(P)-d-my)P$$

$$\frac{dy}{dt} = \begin{pmatrix} \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(P)\ y \end{pmatrix}$$

$$\frac{dz}{dt} = \gamma y - (b(P) - my)z$$

Setting $x_1$=P and $x_2$=(y,z) 'where' denotes transpose yields a structured model where $A_1(x)$=b-d-my and

$$A_2(x) = \begin{pmatrix} \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(p) & 0 \\ \gamma & my\text{-}b(p) \end{pmatrix}$$

Since {0,(*K*,0,0)}is a Morse decomposition of the boundary dynamics Theorem 3 and Remark 1 imply that this model is robustly permanent if and only if b(0)>d and $\beta\frac{K}{\epsilon+K} > \gamma + m + d$. In particular, for $\epsilon > 0$ sufficiently small, one requires that the basic reproductive number $\frac{\beta}{r+m+d}$ is greater then one.

Remarkably this same criterion determines robust permanence of model with significantly more complicated boundary dynamics. For example, suppose the focal population is a predator species. If the prey has abundance *N* and exhibits logistic dynamics f(*N*)is the per-capita predator consumption rate of the prey and b(N)is the per-capita reproductive rate of the predator then the dynamics become

$$\frac{dN}{dt} = rN(1 - N/K - f(N)P$$

$$\frac{dP}{dt} = (b(N) - d - my)P$$

$$\frac{dy}{dt} = \begin{pmatrix} \beta(1 - y - z)\frac{P}{\epsilon+P} - r - m(1-y) - b(N)\ y \end{pmatrix}$$

$$\frac{dz}{dt} = \gamma y - (b(N) - my)z$$

Where r>0 and *K*>0, and *f*(*N*),b(N)satisfy f(0)=b(0)Define $x_1$=*N*, $x_2$=*p*, $x_3$=*(y,z)*′ Let us assume that the disease-free community (i.e.y=z=0)is permanent (i.e.b(K)>d) and let A be the global attractor in the interior of the P-N plane Under these assumption a Morse decomposition of the boundary dynamics is given by $(x_1, x_2, x_3) = (0,0,0)$, $(x_1, x_2, x_3) = (K, 0,0)$and A. Since the equilibria (0,0,0) and (k,0,0) are unsaturated, it remains to characterize $\curlywedge_3(\mu)$for any invariant measure μ supported by A. Proposition 3 and Remark 1 imply

$$\curlywedge_2(\mu) = \int_c (bx_1)\text{-}d)d\mu(x) = 0$$

*Hence* $= \int_c bx_1)d\mu(x) = d$ *and*

$$\curlywedge_3(\mu) = \beta \int \frac{x_2}{\epsilon + x_2} d\mu(x) - r - m - d$$

**Theorem 3**: and Remark 1 implies that this structured model is robustly permanent if $\frac{\beta}{r+m+d} > 1 and\ \epsilon > 0$ is sufficiently small.

6.2. Gene networks. The repressilator is an oscillatory gene network based on three (or more generally an odd number of transcriptional repressors. Such a system has been genetically engineered in E-coli by Elowitz and Leibler [47] Mathematical model go back to [48.49] they all involve concentrations of proteins and mRNs. A modified model with auto-activation was suggested in [50]

$$(14)\ x_i = \beta(y_i - x_i)$$

$$(15)\ y_i = aF(x_{)i} \text{-} y_i)$$

With

$$(16)\quad F(x)_i = x_i g\ (x_i, x_i - 1) = \frac{x_i}{1 + x_i + p x_{i-1} + K p\, x_i x_i - 1}$$

Here $x_i$ and $y_i$ are normalized concentrations of proteins and mRN As belonging to gene $i$,p is the strength of repression and k a cooperativity parameter this is a structured system of type(2)

$$(17)\ \begin{pmatrix} x_i \\ y_i \end{pmatrix} dot = \begin{pmatrix} -\beta & \beta \\ ag(x_i,x_i)-1 & \end{pmatrix} \begin{pmatrix} x_i \\ y_i \end{pmatrix}$$

In [50,Thm 3] it was shown that for n=3,a>1 and p>1,system(14)has a heteroclinic cycle connecting the 3 single gene equilibria (similar to the rock –scissors-paper dynamics of section5.2)Moreover

- the system is permanent i.e. the boundary of $R^6$+is repelling, if $\lambda + \mu > 0$
- this heteroclinic cycle is asymptotically stable if $\lambda + \mu < 0$

where $\lambda > 0$ and $\mu < 0$ are the invasion rates of the two missing genes i.e. the leading eignvalues of the 2X2 matrix in (17)evaluated at the appropriate single gene equilibria (for explicit expressions see[50,eqs.(127,128)].)

Theorem 3 implies the permanence condition implies the stronger conclusion of robust permanence A similar result holds for n odd, but the heteroclinic cycle is then actually a heteroclinic network.

Note: For references please see the original paper.

One of the examples for which we apply the model are as follows: **Progress in Particle and Nuclear Physics Volume 88, May 2016, Pages 1–28 Chiral magnetic and vortical effects in high-energy nuclear collisions—A status report D.E. Kharzeeva, b, J. Liaoc, d, , , S.A. Voloshine, G. Wangf**

The interplay of quantum anomalies with magnetic field and vorticity results in a variety of novel non-dissipative transport phenomena in systems with chiral fermions, including the quark–gluon plasma. Among them is the Chiral Magnetic Effect (CME)—the generation of electric current along an external magnetic field induced by chirality imbalance. Because the chirality imbalance is related to the global topology of gauge fields, the CME current is topologically protected and hence non-dissipative even in the presence of strong interactions. As a result, the CME and related quantum phenomena affect the hydrodynamical and transport behavior of strongly coupled quark–gluon plasma, and can be studied in relativistic heavy ion collisions where strong magnetic fields are created by the colliding ions. Evidence for the CME and related phenomena has been reported by the STAR Collaboration at Relativistic Heavy Ion Collider at BNL, and by the ALICE Collaboration at the Large Hadron Collider at CERN. The goal of the present review is to provide an elementary introduction into the physics of anomalous chiral effects, to describe the current status of experimental studies in heavy ion physics, and to outline the future work, both in experiment and theory, needed to eliminate the existing uncertainties in the interpretation of the data. Keywords Chiral magnetic effect; Chiral vortical effect; Chiral anomaly; Quark-gluon plasma; Heavy ion collisions

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# Sufficient Conditions For The Existence And Linear Stability Of Time Independent Age Distributions In Respect Of Particles, Propositions, Problems, Hegelian Structures, Deleuzean Treatises, Interactions And Transactions In Cosmos And Others

**Abstract:** Gopalaswamy studied sufficient conditions are obtained for the existence and linear stability of time independent age distributions in two species competition with age time lagged density dependent mortality and fertility functions.. Work was published in Bulletin Australian Mathematical Society Volume 25 (1982). We show how the model could be applied in case of finance, econometrics, psycho economics, interactions between particles and transactions in nature. Basically the competition is for resources, notwithstanding such experiments could be conducted in laboratory like a rat kept in front of a cat, such incidents could be represented mathematically, and the corresponding linear stability conditions ca be obtained in case of interactions in cosmos. Take for example the case of a cash credit account of a borrower. Many inward bills waiting for paid, or inland or foreign bills payment all are in line for the debit of the account, which also acts as a resource. In addition there are recurrent cases where in the credit is issued to the cash credit account in case of demand and inland bills. In case of economics, investment has inventories and it has to be utilized towards the end of consummation of output. And this is case of output gobbling part of investment which is inventories. Output has to be sold to the customer and the amount received is diverted partly to savings and consumption and part could be employed towards the consolidation of the system. So this is a case of investment itself acting as a both receiver and payer. It shall eat and also produce the finance. Savings as stated above could be deployed for the further diversification and proliferation and that is towards the end of putting the company on terra firma. Another example in cosmology is the case of black holes guzzling matter. There are lot of black holes and lot of matter. Result of such a guzzling the production of Hawking radiation. Horava gravity is another example. How the conformal field theory could be changed in the case of Horava gravity. Still another case is the case of Vonneumann algebra in the cae of Horava gravity. In this case like cash credit account, the Horava gravity and Von Neumann algebra act as both receivers and takers. Such incidents are legion in cosmology and even in philosophy. We have considered self as individual general ledger, and the debits and credits are raised to the same account thus anointing as both a receiver and payer. We give plethora of such cases and apply the model. Due to spatial constants, same module is used for all the modules and we rewrite to state that the model is applicable to all modules. This is due to spatial restrictions. What is to be done is that each and every parameter is to be changed with a subscript or a superscript and finally it is to be concatenated towards the end of obtention of prediction model or optimizing model, for Duality of motion model for further analysis. Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a

qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

**Introduction**

Let $P_1$ *(t)* and $P_2$ *(t)* denote the total particles, interactions, transactions, etc, which we dilate upon further investigation. Perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. **Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems** sizes (or biomasses) at time $t\geq 0$ of two interacting species living in a common habitat and competing for a common pool of resources. The competition for resources will be implicit in our model similar to that in the two species Lotka- Volterra competition system. Assuming constant sex ratios in the two species we can consider $P_1$ *(t)* and $P_2$ *(t)*to be the Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

population of f; immigration, emigration and internal dispersion in the habitat are assumed to play no significant role in the dynamics of the community.Popoulation we mean Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

We suppose that the two species contain respectively $p_1\,(a, t)\,da$ and $p_2\,(a, t)\,da$ individuals with ages between $a$ and $a + da\ (a\geq 0)$ at time $t$ so that we have

(1.16) $P_1\,(t) = \int_0^\infty \rho_1(a, t)\,da,\ P_2\,(t) = \int_0^\infty \rho_2(a, t)\,da,\ t\geq 0.$

The rates of change of the two species densities are by definition $D\rho_1\,(a, t)$ and $D\rho_2\,(a, t)$ where

(1.17) $D\rho_i\,(a, t) = \limsup_{h\to 0} \frac{\rho_i(a+h,t+h)-\rho_i(a,t)}{h},\ i = 1, 2.$

Assuming the existence of the limits in (1.2) we consider the following time lagged model system

(1.18) $$\begin{cases} D\rho_1\,(a,t) + f_1\,(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_1(a, t) = 0, & a>0,\ t>0 \\ D\rho_2(a,t) + f_2(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_2(a, t) = 0, & \text{(1.4)} \\ \rho_1(0, t) = \int_0^\infty b_1\,(a, P_1\,(t), P_2\,(t-\tau))\rho_2(a, t)da\,, & t>0, \\ \rho_2(0, t) = \int_0^\infty b_2\,(a, P_1\,(t), P_2\,(t-\tau))\rho_1(a, t)da\,, & \end{cases}$$

(1.5) $\rho_i\,(a, s) = \varphi_i\,(a, s),\ a.0,\ s\epsilon\,[-\tau, 0],\qquad i = 1, 2,$

(1.6) $P_i(s) = \int_0^\infty \varphi_i\,(a, s)\,da,\ s\in\,[-\tau, 0],$

Where $\tau$some is fixed nonnegative constant; $f_1$ and $f_2$ denote the age and density dependent mortality rates with time lags in interspecific interactions, $b_1$ and $b_2$ denote the age and density dependent fertility functions again having time lags in the interspecific interactions; $\varphi_1$ and $\varphi_2$ denote the initial age distributions needed for the formulation of the model. Note we mean the population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

The model system (1.1)-(1.6) will be meaningful only if $\varphi_{1,}$ $\varphi_2$, $f_1$, $f_2$, $b_1$, $b2$ are nonnegative functions of the respective arguments. Also since the total initial populations have o be finite, $\varphi_1(\cdot, s)$ should belong to $L_1(R^+)$ for each $s\epsilon$ $[-\tau, 0]$. That $f_i$ and $b_i$ $(i = 1, 2)$ are independent of t explicitly and this assumption can be interpreted to represent the temporal uniformity of the environment.

One of the fundamental questions for (1.1)-(1.6) is the following; under what conditions on $f_{i,}$ $b_i$ and $\tau$ does the system (1.1)-(1.4) have time independent nontrivial solutions, and when such solutions, and when such solutions exist, are they stable with respect to some suitable stability criterion? When such stable time invariant solutions exist for (1.1)-(1.4) we say that the two types has a stable age distribution. Existence of a stable age distribution is the analogue of the existence of a stable steady state for the corresponding age independent system. Types we mean types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

In this article we investigate the above question for the case of a competitive interaction with some additional constraints on the vital rates $f_i$ and $b_i$ $(i = 1, 2)$. Age dependent population systems without time delays have been considered by several authors (Gurt in and MacCamy [3, 4, 5], Gurt in and Levine [6], Haimovici [7] and Rotenberg [8]). The question of the relation between an age independent system and a corresponding age dependent system has been considered by Gurt in and MacCamy [4] and Gurt in and Levine [6] who have established an asymptotic ($as$ $t\rightarrow\infty$ ) relation between such models by constructing a higher dimensional 1 umped parameter system (in terms of ordinary differential equations) to represent the age dependent distributed parameter system. Assuming the existence of stationary age distributions Haimovici [7] considers their stability in a system of two interacting populations explicitly taking into consideration the dynamical nature of the encironments resources and pollution.

In this article we add another realistic feature namely time lags in the age dependent models and obtain sufficient conditions for the existence of stationary age distributions in (1.1)-(1.4) and show that the time lags in interspecific

interactions have no effect on the linear stability of the age distributions although the decay rates of perturbations will depend on the magnitudes of the time delays. Such "harmless" natures of time lags in interspecific interactions have been noted by the author for the age independent models in [1, 2].

## 6. Existence and uniqueness of solutions

We first narrate our assumptions on the vital rates of $f_i$ and $b_i$ $(i = 1, 2)$ so that the system (1.1)-(1.6) will denote a competition system. Let $BC\ (0, \infty)$ denote the linear space of bounded continuous real functions on $(0, \infty)$ with the norm defined by

$$||f|| = \{max\ |\ f(t)\ |,\ t \in (0, \infty)\}.$$

Let $K_1$ be the cone of nonnegative functions of $BC\ [0, \infty]$. Let $C\ [-\tau, 0]$ denote the space of real continuous functions with a norm defined by $||g|| = \{\ max\ |g(s)\ |,\ s \in [-\tau, 0]\}$; let $K_2$ be the nonnegative cone of functions in $C\ [-\tau, 0]$ . We present our first set of assumptions on $f_i$, $b_i$ and $\varphi_i$ for $i = 1, 2$.

$A_1$. $f_1 \in C\ (R^+ \times K_1 \times K_2, R^+)$, $f_1 \geq d_1^*$ on $R^+ \times K_1 \times K_2$

$f_2 \in C\ (R^+ \times K_2 \times K_1, R^+)$, $f_2 \geq d_2^*$ on $R^+ \times K_2 \times K_1$ ,

( $d_1^*$ , $d_2^*$ are positive constants);

$f_1$ and $f_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ ; that is, *for* $i = 1, 2$,

$$|f_i\ (a, P_1^{'}\ , P_2^{'}) - f_i(a, P_1^{"}\ , P_2^{"})| \leq d_i\ [||P_1^{'} - P_1^{"}, P_2^{'} - P_2^{"}||]$$

($d_1$, $d_2$ are positive constants),

$$\partial f_1/\ \partial P_1, \partial f_1/\ \partial P_2 \in C(R^+ \times K_1 \times K_2, R^+)\ ,$$

$$\partial f_2/\ \partial P_1, \partial f_2/\ \partial P_2 \in C(R^+ \times K_2 \times K_1, R^+),$$

$A_2$. $b_1 \in C\ (R^+ K_1 \times\ K_2, R^+)$, $b_1 \leq \beta_1^* on\ R^+ \times K_1 \times K_2$ ,

$b_2 \in C\ (R^+ K_2 \times\ K_1, R^+)$, $b_2 \leq \beta_2^* on\ R^+ \times K_2 \times K_1$ ,

*($\beta_1^*$ , $\beta_2^*$ are positive constants);*

$b_1$ , $b_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ *; that is for* $i = 1, 2$,

$$|b_i\ (a, P_1^{'}, P_2^{'}) - b_i\ (a, P_1^{"}, P_2^{"})| \leq \beta_i\ [||P_1^{'} - P_1^{"}|| + ||P_2^{'} - P_2^{"}]$$

$B_1$ , $B_2$ Are positive constants),

$b_i(\circ,(P_1,P_2)\epsilon L^{\infty}(R^+,R^+)$ For fixed $P_1,P_2,i=1,2,$

$$\frac{\partial b_1}{\partial P_2},\frac{\partial b_1}{\partial P_2}\epsilon\ C\ R^{+}\text{x}\ k_{\ 1}\text{x}\ k_{\ 2},\text{R}^{-}),$$

$$\frac{\partial b_2}{\partial P_1},\frac{\partial b_2}{\partial P_2}\epsilon\ C\ R^{+}\text{x}\ k_{\ 2}\text{x}\ k_{\ 1},\text{R}^{-}),$$

$A_3.\ \varphi_i\ (\cdot\ ,\text{s})\in \text{L}^{\ 2}\ (\text{R}^{+},\text{R}^{+})\ \cap \text{C}(\text{R}^{+},\text{R}^{+})$ for each $s\in\ [-\tau,0],$

$i,\ =1,\ 2,$

By a solution of $(1.1)-(1.6)$ on $[0,\text{T}]$, we mean

$\text{P}_i\text{: R}^{+}\ \square[-\tau,\text{T}]\rightarrow \text{R}^{+}\ \ (i,\ =1,\ 2,)$ With the following properties:

$S_1.$ $\text{P}_i(\cdot\ ,t)\in \text{L}^{\ 1}\ (\text{R}^{+},\text{R}^{+})$ for each $t\ \in[-\tau,\text{T}];$

$S_2.$ $\text{P}_i(a,\cdot\ )\in \text{C}\ ([-\tau,\text{T}],\text{R}^{+})$ for each $a\ \in \text{R}^{+}$;

$S_3.$ $\text{DP}_i$ Exists along the characteristics $t-a=$ constant

On $\text{R}^{+}\ \square\{\text{R}^{+}\cap\ [0,\text{T}]\}$ And is continuous $\ (i,\ =1,\ 2,);$

$S_4.$ $\text{P}_i$ Satisfies the system (1.1)- (1.6) for

$(a,t,)\ \in \text{R}^{+}\ \square[-\tau,\text{T}]$ .

Let us first convert the system (1.1)- (1.6) into an equivalent system of integral equations. For brevity we denote by P the pair ( $\text{P}_1,\text{P}_2$) If we let

$$\text{P}_i(a+x,t+x)=\ \ \text{P}_i(x),(i,\ =1,\ 2,)\ ,$$

Then (1.3) considered at $(1.3)$ considered at $(a+x,\ t+x)$ becomes

$$\frac{dp_{\ i}(\text{x})}{dx}+f_i(a+x,\ P(t+x))\ \ P_i(x)=0$$

Which has a solution of the form?

$$P_i(x)=P_i(0)\ \exp\left[-\int_0^{\infty} f_{\ i}\big(a+s,P(t+S)\big)ds\right]$$

And hence

$$(2.1)\ p_{\ 1}\ (a+x,t+x)=p_{\ 1}\ (a,t)\ exp\big[-\int_o^x\ f_{\ 1}\ \big(a+s,p(t+s)\big)ds\big]$$

For $a \geq t$ *we choose (a,t)=(a-t,0)* in (1.3) and *x=t* in (2.1) so that

(2.2) $p_1(a,t) = \varphi_1(a-t,0)\exp\left[-\int_o^t f_1\left(a-t+s, p_1(s)\right)ds\right]$

$$O \leq t \leq a, i = 1,2,$$

With

$$P_1(s)=(p_1(s),p_2(s-t)),$$

$$P_2(s)=(p_1(s-t),p_2(s)),$$

For $0 \leq a < t$ we choose *(a,t)=(0,t-a)in (1.3)and x=a in (2.1)* so that

(2.3) $p_i(a,t)= B_i(t-a)exp[-\int_o^a f_1(s, p_i(t-a+s))ds]$

$$t > a \geq, i = 1,2,$$

Where

(2.4) $B_i(t) = p_i(0, t), i = 1,2, t > 0.$

Define $M_i and\ L_i\ (i = 1,2)$ as follows:

$$M_i(a,t, p_i(t)) = exp\left[-\int_o^{t-a} f_1(s, p_i(a-a+s))ds\right]$$

(2.5) $L_i\left(a,t,\ p_i(t)\right) = \exp\left[-\int_o^t f_1(\text{a+s},\ p_i(s))ds\right]$

It will now follow from (1.1),(1.3),(2.2)-2.5)that

(2.6) $p_i(t) = \int_o^t B_1(a)\, M_1(a, t, p_i(t))da + \int_o^\infty \varphi_i(a, 0)\ L_i(a,t,p_i(t)da,$

$$t > 0, i = 1,2,$$

(2.7) $p_i(s) = \Phi_i(s) = \int_o^\infty \varphi_i(a, s)da, s \in [-\tau, 0], i = 1,2,$

(2.8) $\beta_i(t) = \int_0^t \beta_i(a) b_i(t-a, P_i(t)) M_i.\left(a,t,P_i(t)\right)da$

$$+\int_0^\infty b_i\left(a+t, P_i(t)\right)\varphi_i(a,0)L_i\left(a,t,P_i(t)\right)da$$

$$i = 1, 2, t > 0$$

the equivalence of (2.6)-(2.8) with (1.1)-(1.6) is established by the following whose proof is identical to a similar result of Gurtin and MacCamy [4]; hence we omit the details of the proof of the following.

THEOREM 1 : *Suppose the system* (1.1)-(1.6) *satisfies the conditions* $A_1, A_2, A_3$ . *If* $(P_1, P_2)$ *is a solution of*

*(1.1)-(1.6) then the total population* $P_1, P_2$ *and the birth rates* $B_1, B_2$ *satisfy the integral equation (2.6)-(2.8).*

*conversely if* $P_1, P_2, B_1, B_2$ *are nonnegative continuous solution of (2.6)-(2.8) on [ 0, T] and if* $P_1, P_2$ *are*

*defined by (2.2)-(2.4) on* $R^+$ x [0, T] *then such* $P_1, P_2$ *provide a solution of (1.1)-(1.6) on* $R^+$ x [0, T] .

The following a *priori* estimates are useful to prove our existence theorem below. It will immediately follow from the bounds for $f_i$ and $b_i$ that

$$M_i(a,t,P) \leq \exp[-d_i^*(t-a)] \leq 1, if \quad t \geq a,$$

*(2.9)*

$$L_i(a,t,P) \leq \exp[-d_i^* t]$$

From (2.7)-(2.8) we derive that

$$\beta_i(t) \leq \int_o^t \beta_i^* \beta_i(a) \exp[-d_i^*(t-a)]\, da + \beta_i^* \Phi_i(0) \exp[-d_i^* t], \quad i = 1,2,$$

And hence by Gronwall's inequality,

(2.10) $$\beta_i(t) \leq \beta_i^* \Phi_i(0) \exp[\delta_i t], \ \delta_i = \beta_i^* - d_i^*$$

(2.5), (2.6) and (2.10) lead to

(2.11) $$|P_i(t) - 0_i(0)| \leq \beta(0) exp[\delta_i a - d_i^*(t-a)]\, da$$

$$+\int_0^\infty \varphi_i(a,0)|1 - exp(-d_i^* t)|da$$

$$\leq \beta_i^* \Phi_i(0) exp(\beta_i^* t)\left\{\frac{1-exp[-(d_i^*+\beta_i^*)t]}{d_i^*+\beta_i^*}\right\} + \beta(0) d_i^* t$$

$$\leq [\beta_i^* \beta(0) exp(\beta_i^* t) \beta(0) d_i^*] t \quad for \ t > 0.$$

We can now prove the following result on the existence and uniqueness of solution of the integral equations (2.)-(2.8)

**THEOREM 2:** *Assume that the system* (1.1)-(1.6) *satisfies the hypotheses*

$A_1, A_2, A_3$. *Then there exists a unique set of continuous function* $(P_1, P_2, B_1, B_2)$ such *that*

$$P_i : R^+ \to R^+, \qquad B_i : R^+ \to R^+ \ (i = 1,2),$$

*Which satisfy the integral equations* (2.6)-(2.8)

Proof. We will first show the existence of local solution on same (possibly small) interval [0, T] $\subset R^+$ such that

$$P_i \in \mathrm{C}([0,T], R^+), \quad B_i \in \mathrm{C}([0,T], R^+), i = 1,2,$$

Satisfying (2.6)-(2.8) and then show the $P_i$ and $B_i$ can be continued as solution for any finite interval in $R^+$ for some fixed

Positive number $T$ we let

$$C^+[-\tau, T] = \{f \in \mathrm{C}([-\tau, T], R) | f \geq 0\}$$

Let $||.||_T$ denote the supremum norm in $\mathrm{C}^+[-\tau, T]$. Then for any

$(x, y) \in \mathrm{C}^+[-\tau, T] \times \mathrm{C}^+[-\tau, T]$ we define

$$||(x,y)|| = ||x||_T + ||y||_T,$$

$$\mathrm{C}_T = \mathrm{C}^+[-\tau, T], \quad \mathrm{C}_T^2 = \mathrm{C}^+[-\tau, T] \times \mathrm{C}^+[-\tau, T].$$

Members of $\mathrm{C}_T$ will be denoted by $Z_t$ with the meaning that

$Z(t+s) = z_t(s)$ *for* $s \in [-\tau, T]$ *and* $t \in [0, T]$. Now for any fixed

$P_t = (P_{1t}, P_{2t}) \in \mathrm{C}_T^2$, (2.8) is a linear system of uncoupled Volterra integral equations in B = $(B_1, B_2)$ and hence (2.8) will

have a unique continuous solution under the assumption $A_1 - A_3$. let us denote such a solution by

$$\mathrm{B}(t) = B^T(P_t), \quad t \in [0, T]$$

Or

$$B_1(t) = B_1^T(P_{1t}, P_{2t}), \quad B_2(t) = B_1^T(P_{1t}, P_{2t}), \ t \in [0, T]$$

If we supply this solution $(B_1, B_2)$ in (2.6) we find that (2.) will be satisfied if and only if $(P_{1t}, P_{2t})$ is a fixed point of the

operator

$$\Pi(P_t) = \{\Pi_1(P_{1t}, P_{2t}),\ \Pi_2(P_{1t}, P_{2t})\}$$

Where $\pi_1$ *and* $\pi_2$ respectively denote the right sides of (2.6) for *i =1,2* .

It is not difficult to see from the nature of $M_i$ , $L_i$ and the hypotheses on $f_i, b_i, \varphi_i$ that

$$B: C_T^2 \to C_T^2 \text{ and } \Pi C_T^2 \to C_T^2$$

(which guarantees the nonnegativity of $P_i$ and $B_i$). We will first show that

There exist positive constants $r^*, t^*$ such that $\pi$ is a mapping of a sphere

$$S_{t*}(\Phi, r^*) = \{ P_t \in (P_{1t}, P_{2t}) \in C_{t*}^2 \mid \|P_t - \Phi_t\|_{t*} \le r^* \}$$

Into itself and also is a contraction where $\Phi_t$ is defined on [-$\tau$, $t^*$]

By the following:

$$\Phi_t = (\Phi_{1t}, \Phi_{2t});$$

(2.12)

$$\Phi_{it} = \begin{cases} \int_0^\infty \varphi_i(a, s)\, da \text{ for } s \in [-\tau, 0]. \\ \int_0^\infty \varphi_i(a, 0)\, da \text{ for } s \in [0, t^*]. \end{cases}$$

We note that similar to (2.11) one gets the estimates

$$|\Pi(P)(t) - \Phi_i(0)| \le \{B_i^* \Phi_i(0) \exp[B_i^* t] + \Phi_i(0) d_i^*\} t$$

Which together with (2.7) and (2.12) lead to

(2.13) $$\|\Pi(P_t) - \Phi_t\|_{t*} \le \left( \sum_{j=1}^2 \|\Phi_{jt}\|_{t*} [B_j^* \exp(B_j^* t) + d_i^*] \quad t^* \right)$$

Thus if we choose,

$$r^* = T \sum_{j=1}^2 \|\Phi_{jt}\|_{t*} [B_j^* \exp(B_j^* T) + d_i^*]$$

Then $\Pi$ is a mapping of $S_{t*}(\Phi, r^*)$ into itself for all $t^* \le T$. we then will have to show that for small $t^* > 0$ , $\Pi$ is a contraction.

Let $P_t^{(1)}$ *and* $P_t^{(2)}$ be arbitrarily chosen in $S_{t*}(\Phi, r^*)$ and consider $|\Pi(P_t^{(1)}) - \Pi(P_t^{(2)})|$ for $t \le t^*$.

Corresponding to the chosen $P_t^{(1)}$ , $P_t^{(2)}$ we let $B^{(1)}(a) = B(P_T^{(1)})$ and $B^{(2)}(a) = B(P_T^{(2)})$ . Then from

(2.6),

$$|\Phi(P^{(1)}(t)) - \Phi(P^{(2)}(t)| \leq \left| \int_0^t B_1^{(1)}(a)\left\{M_1\left(a,t,P_t^{(1)}\right) - M_1\left(a,t,P_t^{(2)}\right)\right\}da \right|$$

$$+ \left| \int_0^t [B_1^{(1)}(a) - B_2^{(2)}(a)]\, M_1(a,t,P_t^{(2)})da \right|$$

$$+ \left| \int_0^t \varphi_1(a,0)\left[L_1\left(a,t,P_t^{(1)}\right) - L_1\left(a,t,P_t^{(2)}\right)\right]da \right|$$

$$= J_1 + J_2 + J_3 \text{ (say)}$$

using the elementary inequality

$$|e^x - e^y| \leq |x-y| \ for \ x,y \leq 0$$

We derive from (2.5) that

$$\left| M_1(a,t,\ p_t^{(1)}) - M_1(a,t,\ p_t^{(2)}) \right|$$

$$\leq \int_o^{t-a} \left| f_1(s,\ p_t^{(1)}(a+s)) - f_1(s,\ p_t^{(2)}(a+s)) \right| ds$$

$$\leq d_1 \| p_t^{(1)} - p_t^{(2)} \|_{t^*} (using\ A_1)$$

*Now,using (2.10),*

$$\left| J_1 \right| \leq \int_o^a \beta_1^* \Phi_1(0)\, e^{\delta 1^a} \left(\int_o^{t-a} d_1 \| P_t^{(1)} - P_t^{(2)} \|_{t^*} ds\right) da$$

$$\leq d_1 \beta_1^* \Phi_1(0)\, t^2 \| P_t^{(1)} - P_t^{(2)} \|_{t^*} e^{\delta 1^t/2}$$

In a similar way

$$\left| J_3 \right| \leq \int_o^\infty \varphi(a,o)\left[\int_o^t \left| f_1(a+s, P_t^{(1)}(s) - f_1(a+s, P_t^{(2)}(s)) \right| ds\right] da$$

$$\leq d_1 \Phi_1(0) \| P_t^{(1)} - P_t^{(2)} \|_t$$

Now we have to estimate $\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right|$ in order to estimate $J_2$, from (2.8),

$$B_1^{(1)}(a) - B_1^{(2)}(a)$$

$$= \int_o^a \left[B_1^{(1)}(\sigma) - B_1^{(2)}(\sigma) b_1\left(a-\sigma, P_t^{(1)}(a)\right) M_1(\sigma, a, P_t^{(1)}(a) d\sigma + \int_o^a \left[B_1^{(2)}(\sigma)\right.\right.$$

$$\mathrm{X}\left[b_1(a-\sigma, P_t^{(1)}(a) M_1(\sigma, a, P_t^{(1)}(a) - b_1\left(a-\sigma, P_t^{(1)}(a) M_1\left(\sigma, a, P_t^{(2)}(a)\right)\right)\right] d\sigma$$

$$+\int_o^\infty \varphi(\sigma, o)\left[b_1\left(\sigma + a, P_t^{(1)}(a)\right) - b_1\left(\sigma + a P_t^{(2)}(a)\right)\right], L_1\left(\sigma, a, P_t^{(1)}(a)\right) d\sigma$$

$$+\int_o^\infty \varphi(\sigma, o) b_1\left(\sigma + a, P_t^{(2)}(a)\right)\left[L_1\left(\sigma, a, P_t^{(1)}(a)\right) - L_1\left(\sigma, a, P_t^{(2)}(a)\right)\right] d\sigma$$

$$= R_1 + R_2 + R_3 + R_4 \text{ (say)}.$$

By the uniform Lipschitz continuity of $b_1$ in $A_2$ we get

$$\left| b_1\left(a+\sigma, P_t^{(1)}(a)\right) - b_1\left(\sigma + a, P_t^{(2)}(a)\right) \right| \leq B_1 \, /\!/ \, P_t^{(1)} - P_t^{(2)} /\!/ \text{ for } a \, \varepsilon [o, t^*]$$

And hence we have

$$|R_3(a) + R_4(a)| \leq \beta_1 \Phi_1(0) || P_t^{(1)} - P_t^{(2)} ||_{t^*} + \Phi_1(0) \beta_1 d_1 || P_t^{(1)} - P_t^{(2)} ||_{t^*}^{*a},$$

$$|R_2(a)| = [\int_0^a [B_1^{(2)}(\sigma)\{b_1(a-\sigma, P_t^{(1)}(a)) - b_1(a-\sigma, P_t^{(2)}(a))\} M_1(\sigma, a, P_t^{(1)}(a))$$

$$+ b_1(a, -\sigma, P_t^{(2)}(a))\{M_1(\sigma, a, P_t^{(1)}(a)) - M_1(a-\sigma, P_t^{(2)}(a))\}] d\sigma]$$

$$+ b_1\left(a - 0, P_t^{(2)}(a)\right)\{M_1(\sigma, a, P_t^{(1)}(a)) - M_1(a-\sigma, P_t^{(2)}(a))\}] d\sigma]$$

$$\leq \int_0^a [\beta_1^* \Phi_1(0) \, \mathrm{e}^{\delta 1^a} \beta_1 \, \| P_t^{(1)} - P_t^{(2)} /\!/_{t^*} + \beta_1 \mathrm{d}_1 \, \| P_t^{(1)} - P_t^{(2)} /\!/_{t^*} (a-o)] do$$

$$\leq \beta_1^* \Phi_1(0) \, \mathrm{e}^{\delta 1^a} \| P_t^{(1)} - P_t^{(2)} /\!/_{t^*} \{\beta_1 a + d_1 a^2 / 2\}$$

Thus

$$\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right|$$

$$\leq \int_o^a \left| B_1^{(1)}(\sigma)\text{-}B_1^{(2)}(\sigma) \right| B_1^{*d}\sigma$$

$$+ \Phi_1(0)/\!/ P_t^{(1)}\text{-}P_t^{(2)}/\!/_{t^*}\left\{ B_1 + B_2\, d_1\, a + B_1^{*} \left( B_2 a + d_1 a^2/2 \right) e^{\delta} 1^{a} \right\}$$

$$\leq \beta_1^{*} \int_o^a \left| B_1^{(1)}(\sigma)\text{-} B_1^{(2)}(\sigma) \right| \text{do+} \left( r^{*} + /\!/ \Phi_t /\!/_{t*} \right) \| \; P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} g_1^{(a)}$$

Where $g_1(a) \rightarrow \beta_1$ as $a \rightarrow 0$. By Gronwall's inequality it will follow that

$$\left| B_1^{(1)}(a)\text{-}B_1^{(2)}(a) \right| \leq \left( r^{*} + /\!/ \Phi_t /\!/_{t*} \right) /\!/ \; P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} \{ g_1(a) + 0(a) \}$$

Using this estimate for $R_1(a)$ we get

$$|R_2(a)| \leq \left| \int_0^a [B_1^{(2)}(\sigma)\, B_1^{(2)}(\sigma),]\; b_1(a)) \text{-} b_1(a\text{-}\sigma,\, P_t^{(1)}(a))\, M_1(\sigma, a, P_t^{(1)}(a)) do \right|$$

$$\left| B_1^{*} \left( r^{*} + /\!/ \Phi_t /\!/_{t*} \right) /\!/ \; P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} 0(a) as\ a \rightarrow 0 \right.$$

*Thus*

$$\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right| \leq (r^{*} + /\!/ \Phi_t /\!/_{t*}) /\!/ \; P_t^{(1)} - P_t^{(2)} /\!/_t o(1) as \quad a \rightarrow o$$

And hence

$$\left| \pi_1(P_t^{(1)} - \pi_1\left(P_t^{(2)}\right) \right| \leq ((r^{*} + /\!/ \Phi_t /\!/_{t*}) /\!/ \; P_t^{(1)} - P_t^{(2)} /\!/_t o(t) as \quad t \rightarrow o$$

Which together with a similar bound for $\left| \pi_2(P_t^{(1)} - \pi_2\left(P_t^{(2)}\right) \right|$ leads to

$$\| \; \pi\,(P_t^{(1)} - \pi\,\left(P_t^{(2)}\right) /\!/ \leq (r^{*} + /\!/ \Phi_t /\!/_{t*}) /\!/ \; P_t^{(1)} - P_t^{(2)} /\!/_{t*} o(t^{*})$$

As $t^{*} \rightarrow o$ showing that if $t^{*}$ is small enough, the mapping $\pi$ on

$s_{t^*}(\Phi, r^{*})$ is a contraction. Hence the system (2.6)-(2,8) has a unique solution for $o \leq t \leq t^{*} for\ t^{*}$ sufficiently small and positive

Now if [0,T] is any finite interval on $R^{+}$, we can consider [0,T] as a finite union of intervals of length less than or equal to $t^{*}$ Since the estimate (2.11)is valid for any finite interval we can extend the solution of (2,6)-(2.8) form $[0, t^{*}]$ to [0,T] and this completes the proof,

3. Stationary age distributions and their linear stability we will now establish the existence of time independent solution

$P_t^{*}, P_t^{*}$ Of (1.1)-(1.6): such solutions satisfy

(3.1) $\frac{dp\,P_2^*(a)}{da}+f_i(\mathrm{a}, P_1^*, P_2^*)\,P_i^*(a) = 0,\ \ i = 1,2, a > 0,$

(3.2) $P_i^* = \int_o^\infty P_i^*(a)da, i = 1,\ 2,$

(3,3) $P_i^*=\int_o^\infty b_i(a,\ P_1^*,\ P_2^*)P_i^*(a)da,\ i = 1.\ \ 2$

All solutions of (3.1)are of the from

(3.4) $P_i^*(a) = P_i^*(o)\exp\left[-\int_o^\infty f_i(s, P_1^*, P_2^*)ds\right], i = 1,\quad 2, a \geq 0,$

Which together with (3.3) lead to

$1=\int_o^\infty b_1(a, P_1^*, P_2^*)\exp\left[-\int_o^a f_1(s, P_1^*, P_2^*)ds\right]da,$

(3.5) $1=\int_o^\infty b_2(a, P_1^*, P_2^*)\exp\left[-\int_o^a f_2(s, P_1^*, P_2^*)ds\right]da$

Thus the existence of stationary solutions of (1.1)-(1.6) reduces to the existence of a pair of positive constants $P_1^*, P_2^*$ satisfying (3.5). We can now prove the following,

THEOREM 3. Assume that $f_1$ and $b_i$ satisfy ( $A_1$ ) and ( $A_2$ )and let

(3.6) $F_{i.}\ (P_1, P_2) = \int_o^\infty b_i(a, P_1, P_2)\exp\left[-\int_o^\infty f_i(s, P_1, P_2)ds\right]da,$

$i$ =1, 2, $P_1 \geq 0$, $P_2 \geq 0$,

**And suppose the following hold:**

(3.7) $A_{4.}$ $F_{i.}\ (0,\ 0) > 1,\ i = 1,2$:

$A_{5.}$ *there exists a positive constant c for which*

$F_1$(C,0)< $F_{2.}$(C,0),

(3.8) $F_1$(0.C)< $F_{2.}$(0,C),

Then there exists a unique pair $(P_1^*, P_2^*)$ of real numbers $P_1^* > o$,

$P_2^* > 0$ such that

(3,9) $F_1(P_1^*, P_2^*) = 1 = F_2(P_1^*, P_2^*)$

And the unique nonegative solution of (3.1)-(3.3) is given by

(3.10) $P_2^*(a) = [P_2^* \exp[-f_i(s, P_1^*, P_2^*)ds])/(\int_o^\infty \exp[-\int_o^a f_i(s, P_1^*, P_2^*)ds]da)$,

$i$ =1, 2, $a \geq 0$,

Proof Consider the elements of the surfaces defined by

Z= $F_1(P_1, P_2), z = F_2(P_1, P_2), P_1 \geq 0, P_2 \geq 0$,

In the ($P_1$,$P_2$,z) space. By hypothesis ($A_1$) we have $\partial F_i / \partial P_j < 0$ *and* , by ($A_2$), $\partial b_i / \partial P_j \leq 0$ for all

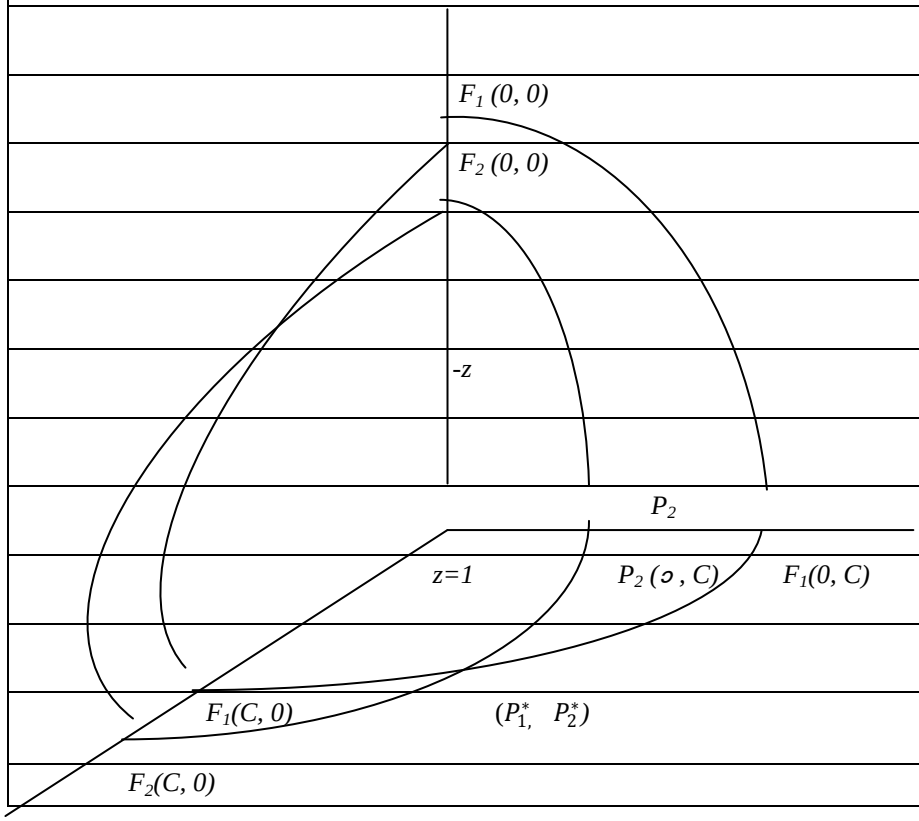
$P_1 \geq 0, P_2 \geq 0$, $i$ =1, 2,

$j$ =1, 2, Hence by ($A_4$) the intersections of these surfaces with the plane z=1 define two curves on the z=1 plane connecting the I ines {z=1, $P_1 \geq 0, P_2 \geq 0,$} and {z= $1, P_2 \geq 0, P_1 \geq 0$, These two curves are defined by

$F_1(P_1, P_2 = 1$ *and* $F_2(P_1, P_2) = 1, F_1(P_1, P_2 \geq 0$,

Which by ($A_s$) will intersect at a unique point say $((P_1^*, P_2^*)$ with $(P_1^* > 0, P_2^* > 0$ (see figure). The uniquess of the point $((P_1^*, P_2^*)$ *is a* consequence of the smoothness of the surfaces $z = F_1(P_1, P_2)$ and $z = F_2(P_1, P_2)$

Following is a hypothetical case for Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates attributable or ascribable in the transactions and reactionary potentialities.



(The figure corresponds to the case $F_1(0,0) > F_2(0,0) > 0$. The other possibility $1 < F_1(0,0)$ is treated similarly) It will now

Follow that for such a pair $(p^*_1, p^*_2)$ the unique nonnegative solution of (3.1) - (3.3) is given by (3.10).

Now to examine the linear stability of the stationary age distributions $p^*_1, p^*_2$, we let

(3.11)
$$\begin{cases} p_i(a,t) = p^*_i(a) + u_i(a,t) \\ p_i(t) = p^*_i + p_i(t), \quad i = 1, 2, \quad a > 0, \quad t > 0, \\ p_i(t) = \int_0^\infty u_i(a,t)da, \end{cases}$$

In (1.1)-(1.5) and derive the following variational system after neglecting the nonlinear terms in the perturbations $\boldsymbol{u}_i$ and $\boldsymbol{p}_i$ $(i = 1, 2,)$ (such a procedure of linear stability analysis can be justified As has been done in Gurtin and MacCamy [4].)

(3.12)
$$\begin{cases} Du_i(a,t) = -f_i\left(a, p^*_1, p^*_2\right) u_i(a,t) - p^*_i(a)\left\{\sum_{i=1}^{2} \frac{\partial f_i}{\partial p^*_i} p_i(t)\right\}, \\ u_i(0,t) = \int_0^\infty b_i\left(a, p^*_1, p^*_2\right) u_i(a,t)da + \int_0^\infty p^*_i(a)\left\{\sum_{i=1}^{2} \frac{\partial b_i}{\partial p^*_i} p_i(t)\right\} da \end{cases}$$

Where $i = 1,2$ and (3.13)
$$p_{11}(t) = p_1(t), \quad p_{12}(t) = p_2(t-\tau), \quad p_{21}(t) = p_1(t-\tau), \quad p_{22}(t) = p_2(t)$$

To consider the asymptotic bahaviour as $t \to \infty$ as of solutions of (3.12) we let

(3.14)
$$\begin{cases} u_i(a,t) = u_i(a) \exp[\lambda t], \quad u_i(a) \to 0 \text{ as } a \to \infty, \\ p_i(t) = p^*_i \exp[\lambda t], \\ \qquad\qquad i = 1, 2, \quad a > 0, \quad t > 0, \\ p^*_i = \int_0^\infty u_i(a)da, \end{cases}$$

Using (3.14) in (3.12) - (3.13),

$$\frac{du_i}{da} + \lambda u_i = -f_i\left(a, p^*_1, p^*_2\right) u_i(a) - p^*_i(a)\left\{\sum_{j=1}^{2} \frac{\partial f_i}{\partial p^*_j} p^*_i\right\},$$

(3.15)

$$\phi_i(0) = \int_0^\infty b_i\left(a, p_1^*, p_2^*\right)\phi_i(a)da + \int_0^\infty p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial b_i}{\partial p_j^*}\, p_i^*\right\}da,$$

Where $i = 1, 2,$ and

$$p_{11}^* = p_1^*, \quad p_{12}^* = p_2^* \exp\{-\lambda\tau\},$$

$$p_{21}^* = p_1^*, \exp\{-\lambda\tau\}, \; p_{22}^* = p_2^*.$$

If we integrate the first of (3.15) and use the fact that $\phi_i(a) \to 0$ as $a \to \infty$ we can derive that $p_1^*$ and $p_2^*$ are governed by

(3.16) $$\lambda p_1^* - \int_0^\infty p_1^*(a)\left[\left(\frac{\partial b_1}{\partial p_j^*} - \frac{\partial f_1}{\partial p_1^*}\right)p_1^* + \left(\frac{\partial b_1}{\partial p_2^*} - \frac{\partial f_1}{\partial p_2^*}\right)p_2^* e^{-\lambda\tau}\right]da$$

$$= \int_0^\infty \left[b_1\left(a, p_1^*, p_2^*\right) - f_1\left(a, p_1^*, p_2^*\right)\right]\phi_1(a)da,$$

(3.17) $$\lambda p_2^* - \int_0^\infty p_2(a)\left[\left(\frac{\partial b_2}{\partial p_1^*} - \frac{\partial f_2}{\partial p_1^*}\right)p_1^* e^{-\lambda\tau} + \left(\frac{\partial b_2}{\partial p_2^*} - \frac{\partial f_2}{\partial p_2^*}\right)p_2^*\right]da$$

$$= \int_0^\infty \left[b_2\left(a, p_1^*, p_2^*\right) - f_2\left(a, p_1^*, p_2^*\right)\right]\phi_2(a)da,$$

From the first of (3.15) we get

(3.18) $$\phi_i(a) = \left\{\phi_i(0)p_i^*(a)/p_i^*(0)\exp[\lambda a]\right\}$$

$$- p_i^*(a)\exp[-\lambda a]\int_0^a \left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\, p_i^*\right\}e^{\lambda s}ds, \quad i = 1, 2,$$

Which by the second of (3.15) becomes,

(3.19) $$\phi_i(a)\left[p_i^*(a)/p_i^*(0)c_i\exp(\lambda a)\right]\left[\int_0^a p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\, p_i^*\right\}da\right.$$

$$\left. - \int_0^\infty b_i\left(a, p_1^*, p_2^*\right)p_i^*(a)\, e^{-\lambda a}\left\{\int_0^a\left(\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\, p_i^*\right)e^{\lambda s}ds\right\}da\right]$$

$$- p_i^*(a)\, e^{-\lambda a}\int_0^a \left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\, p_i^*\right\}e^{\lambda s}ds, i = 1, 2,$$

where

(3.20) $$c_i = 11 - \frac{1}{p_i^*(0)} \int_0^\infty b_i\left(a, p_1^*, p_2^*\right) p_i^*(a) e^{-\lambda a} da, \quad i = 1, 2,$$

Now supplying $\varphi_1(a)$ and $\varphi_2(a)$ from (3.19) in (3.16)-(3.17) we derive that

$$\left(\lambda - A_{11}\right) p_1^* - A_{12} p_2^* \exp\left(-\lambda\tau\right) = g_{11}(\lambda) p_1^* + g_{12}(\lambda) p_2^* \exp(-\lambda\tau),$$

(3.21)

$$-A_{21} p_1^* \exp\left(-\lambda\tau\right) + \left(\lambda - A_{22}\right) p_2^* = g_{21}(\lambda) p_1^* \exp\left(-\lambda\tau\right) + g_{22}(\lambda) p_2^*$$

Where

(3.22) $$A_{i\varphi} = \int_0^\infty p_i^*(a) \left[\frac{\partial b_i}{\partial p_\varphi^*} - \frac{\partial f_i}{\partial p_\varphi^*}\right] da,$$

(3.23) $$g_{i\varphi}(\lambda) = \left[p_i^*(0) c_i\right]^{-1} \left\{\int_0^\infty \sigma_i(a) p_i^*(a) e^{-\lambda a} da\right\} \varphi \left\{\int_0^\infty \left[p_i^*(a) \frac{\partial b_i}{\partial p_\varphi^*} - b_i\left(a, p_1^*, p_2^*\right) p_i^*(a) e^{-\lambda a} \left(\int_0^a \frac{\partial f_i}{\partial p_\varphi^*} e^{\lambda s} ds\right)\right] da\right\}$$

$$- \int_0^\infty \sigma_i(a) p_i^*(a) e^{-\lambda a} \left(\int_0^a \frac{\partial f_i}{\partial p_\varphi^*} e^{\lambda s} ds\right) da, \quad i, \varphi = 1, 2,$$

(3.24) $$\sigma_i(a) = b_i\left(a, p_1^*, p_2^*\right) - f_i\left(a, p_1^*, p_2^*\right), \quad i, = 1, 2.$$

To solve the linear perturbational system (3.12) it is enough to find the constants $\lambda, p_1^*, p_2^*$ satisfying (3.21) since one can then use (3.19) to find the perturbations from (3.14). A nontrivial solution $\left(p_1^*, p_2^*\right)$ for (3.21) will exist if and only if $\lambda$ is a root of the equation $et \begin{bmatrix} \lambda - \left(A_{11} + g_{11}(\lambda)\right) & -\left(A_{12} + g_{12}\right) e^{-\lambda\tau} \\ -\left(A_{21} + g_{21}\right) e^{-\lambda\tau} & \lambda - \left(A_{22} + g_{22}(\lambda)\right) \end{bmatrix} = 0$

Or equivalently

(3.25) $$\lambda^2 - \lambda\left(A_{11} + A_{22}\right) + A_{11}A_{22} - A_{12}A_{21} \exp\left[-2\lambda\tau\right] = \lambda s_1(\lambda) + s_2(\lambda, \tau)$$

Where $S_1(\lambda) = g_{11}(\lambda) + g_{22}(\lambda)$ ,

$S_2(\lambda, \tau) = \left[A_{21} g_{12}(\lambda) + A_{12} g_{21}(\lambda) + g_{12}(\lambda) g_{21}(\lambda)\right] \exp\left[-2\lambda\tau\right]$

$$-A_{11} g_{22}(\lambda) - A_{22} g_{11}(\lambda) - g_{11}(\lambda) g_{22}(\lambda)$$

The following result establishes the linear local asymptotic Stability of the stationary age distributions $p_1^*(a)$ and $p_2^*(a)$

THEOREM 4 Assume the following:

$A_6.\ 0 < \beta_i^* < d_i^*\ (i,\ =1,\ 2.)\ \ (see(A_1)\ and\ (A_2));$

$A_7.\ \{A_{11} + A_{22} + \| S_1(0) \|\}^2 < 4\{(A_{11}A_{22} - A_{12}A_{21}) - \| S_2(0) \|\},$

$$\| S_1(0) \| = |g_{11}(0)| + |g_{22}(0)|,$$

$$\| S_2(0) \| = |A_{11}\, g_{12}(0)| + |A_{12}\, g_{21}(0)| + |g_{12}(0)\, g_{21}(0)|$$

$$+ |A_{11}\, g_{22}(0)| + |A_{22}\, g_{11}(0)| + |g_{11}(0)\, g_{22}(0)|$$

Then all the roots of (3.25) have negative real parts and hence the

perturbations $u_i\ (a,\ t) = \phi(a) \exp[\lambda t] \to 0\ as\ \ t \to \infty$ for $i = 1, 2\ and\ a \geq 0$

Proof . Define $F_1(\lambda, \tau)$ and $F_2(\lambda, \tau)$ as follows:

$F_1(\lambda, \tau) = \lambda^2 - \lambda(A_{11} + A_{22}) + A_{11}A_{22} - A_{12}A_{21} \exp[-2\lambda\tau];$

$F_2(\lambda, \tau) = \lambda S_1(\lambda) + S_2(\lambda, \tau),$

Using $(A_6)$ one can show that for $\lambda$ with re $(\lambda) \geq 0$, $c_1$ and $c_2$ are positive and bounded away from zero, Now since $A_{ij} > 0$ and $\tau \geq 0$ we find from the nature of the dependence of $g_{ij}$ on $\lambda$ that for all $\lambda\ \ \mu \pm iw$ with $\mu \geq 0$ and $w \geq 0$ we have

(3.26) $$\left| F_1(\lambda, \tau) - P_2(\lambda, \tau) \right|$$

$$\geq \left| F_1(\lambda, 0) | - F_2(\lambda, \tau) \right|$$

$$\geq |\lambda|^2 - |\lambda|\{A_{11} + A_{22} + \| S_1(0) \|\} + \{A_{11}A_{22} - A_{12}A_{21} - \| S_2(0, \tau) \|\}$$

$$\geq \left[(|\lambda| - P)^2 + (q)\right]^2$$

Where

$p = \{A_{11} + A_{22} + \| S_1(0) \|\}/2,$

$q = \left\{4 | (A_{11}A_{22} - A_{12}A_{21}) - \| S_2(0, T) \| | - | A_{11} + A_{22} + \| S_1(0) \| |^2\right\}/2,$

It will now follow, from (A) and (3.26), $F_1\ (\lambda, \tau) = F_2(\lambda, \tau)$ connot have roots with zero or positive real parts and this completes the proof.

In conclusion we remark that the condition (3.7) means that the net creation rates of each system is greater than utility; a condition of this type is well known in age wise classification of particles, atoms, molecules, cosmological phenomenon. (3.8) can be interpreted to mean that the intraspecific inhibition is higher than that of the

interspecific interaction. The conditions of Theorem 4 are analytical and do not lend themselves for any worthwhile interpretation in terms of the model parameters. So also an interpretation can be recorded as a remark in respect of destruction rates of particles, properties, elements, or projections.

Note: The model is applicable to each and every module. As stated in the foregoing, there are possibilities of situation where an account itself be a receiver and it is also the one to which debit is raises. Such examples are legion in all the fields of finance, economics, political science and physics and chemistry, and of course mathematics. Classification based on time is the preeminently and predominantly suited premise for such model. We repeat that the model is to be read with each module We will expatiate the importance and implicational ramifications as we proceed with the work. Assortment-blend, collection-combinations, combo-gallimaufry, goulash-hash, hodgepodge-medley, mishmash-mixture, motley-pastiche patchwork-salmagundi, soup-stew, mixed bag-mélange, olio- olio podrida, assortment variety, Array choice, collection combination, combo diversity, hodgepodge jumble, kind medley, miscellany mishmash, mixed bag mixture, mélange potpourri, selection sort, brew concoction, beverage blend, broth compound, distillation drink, fermentation hash, hodgepodge infusion, instillation liquor, mélange miscellany, mishmash mixture, mélange potpourri, borscht bouillon, bowl brew, chowder concoction, consommé decoction, gumbo hodge-podge, olio porridge, potage potpourri, pottage puree, splash stock, vichyssoise hodgepodge, mixture- mess, collection combination, goulash hash, jumble medley, miscellany mishmash, mixed bag mélange, olio patchwork, potpourri salmagundi, jumble clutter, disarrangement disarray, disorder farrago, gallimaufry garbage, goulash hash, litter medley, mess miscellany, mishmash mixture, muddle mélange, olio pastiche, patchwork potpourri, salmagundi scramble, shuffle snarl, tangle tumble, medley miscellany, assortment brew, collection combo, conglomeration farrago, hodgepodge jumble, melee mingling, mishmash mixed bag, mixture mélange, pasticcio pastiche, patchwork potpourri, salmagundi variety, mélange mixture, assortment combo, farrago gallimaufry, pasticcio pastiche, patchwork potpourri, salmagundi soup, stew miscellany, accumulation aggregation, anthology assortment, brew collectanea, combination combo, compilation conglomeration, accumulation diversity, farrago gallimaufry, melee mishmash, mix mixed bag, mixture muddle, mélange odds and ends, olio pasticcio, pastiche patchwork, potpourri salad, salmagundi smorgasbord, stew variety, mishmash jumble, medley mess, miscellany mixed bag, mixture mélange, olio patchwork, potpourri salmagundi, mixed bag array, assortment choice, collection combination, combo conglomerate, diversity hodgepodge, jumble medley, mélange miscellany, mishmash mixture, montage mélange, odds and ends, patchwork potpourri, smorgasbord sundries, mix/mixture assortment, combination, admixture adulteration, alloy amalgam, amalgamation assimilation, association batter, blend brew, combine combo, commixture composite, compound concoction, confection conglomeration, cross crossing dough fusion, goulash grab bag, hodgepodge hybrid, hybridization incorporation, infiltration interfusion, jumble mash, medley merger, mingling miscellany, mishmash mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, odds and ends assortment, bits and pieces, debris etcetera, hodgepodge jumble, knickknacks leavings, leftovers litter, medley melee, miscellaneous paraphernalia, miscellany mishmash, mixed bag motley, mélange notions, novelties oddments, olio particles, patchwork potpourri, rags refuse, remainder remains, remnants rest, rummage scraps, sundries sundry items, compilation mishmash, paste-up patchwork, potpourri reappropriation, reproduction synthesis, conglomeration cross, crossing dough, fusion goulash, interfusion jumble, mash medley, merger mingling, miscellany mishmash, mix mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, colluviums gallimaufry, hash mixture, pasticcio pastiche, conglomeration cross section, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multiformity multiplicity, mélange potpourri, range shift, variegation variousness, divergencies diverseness, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multifariousness multiformity,

mélange potpourri, range shift, soup stew, variance variation, variegation variousness, oddment odds and ends, assortment bits, polymorphism assortment, change collection, combo conglomeration, cross section departure, discrepancy disparateness, divergencies diverseness, diversification diversity, fluctuation heterogeneity, heterogeneousness incongruity, intermixture many-sidedness, medley miscellaneousness, miscellany mishmash, mixed bag mixture, notions novelties, oddment oddments, olio particles, patchwork potpourri of the particles, interactions, accounts, transactions in the cosmos. The wisdom and understanding of our own divine essence is like a beaconing begorra, bureaucratic burdock, befuddling binary and digital digamma, that invites us inward toward ever deeper connection with all around us. Yet rather than look within, most in this world prefer to look outward for divine wisdom, thus sustaining a dependence upon a vast hierarchy that stretches between the individual and the Divine. In all our wanderings away from the Divine, humankind has obscured, obfuscated, obliterated, adjured, abnegated and abjured its most compelling features through a persistent belief systems in limitations arising out of the controlling mechanisms of this hierarchy of sub ordinates, superordinates and coordinates. Has my consciousness changed? Why that doubt? Twist and shake, ponderous novelty themes, pathos, poignancy in writing in different styles, accord-coetaneousness, coevality-coincidence, concurrence-conformity, conjunction-contemporaneousness, harmony, order, peace, simultaneousness- synchronicity exciting optimism of youth, power obmbudsman,pulsations and throbbing of pop, hard luck stories of people, modernity and mischief, grimalkins of rurality, wobbling boats of lovers, saxophone in mouths all point to newer value orientations. The Divine dances and the “Grand design “is perpetrated outside of the confines of any hierarchical structures, be it feudal, organizational. It is complete within itself, and has a singular purpose of demonstrating the collective potential of all life within the universe. It is the archetype of perfection. It is the standard bearer of each soul's innate design and ultimate destiny. The essence of the Divine is far beyond mental conception, yet humanity's tendency is to resort to the limiting language of the hierarchical paradigm acceding- accession, accommodation- accord, accordance, acknowledging, adjustment-affiliation, affinity-alliance, amity-approving, arbitration-arrangement, assenting, authorizing, bargaining-compatibility, compliance--compromise, concert-concession, conformity-congruity, consistency-correspondence, endorsing, granting-harmony in conformity to pen chance ,predilections of the protagonist to define and understand it. Searching for Connection and Wholeness calls for not lumping elements like me made of clay feet and weak knees, owing allegiance to the fringe, without courante and courage of conviction and will for vindication, visual restlessness and exaggerated narratives of self to its own, exciting potentialities available, terrestrial passion, and skew of scams all bear an infallible demonstration for new order of consciousness. Theorists have posited a virtual smorgasbord of objects to account for dark matter, although many of them have fallen prey to observational constraints. As leading possible candidates for baryonic dark matter, there are black holes (large and small), brown dwarfs (stars too cold and faint to radiate), sun-size MACHOs, cold gas, dark galaxies and dark clusters, to name only a few. The range of particles that could constitute nonbaryonic dark matter is limited only slightly by theorists' imaginations. The particles include photinos, neutrinos, gravitinos, axions and magnetic monopoles, among many others. Of these, researchers have detected only neutrinos--and whether neutrinos have any mass remains unknown. Experiments are under way to detect other exotic particles. If they exist, and if one has a mass in the correct range, then that particle might pervade the universe and constitute dark matter. But these are very large "ifs." To a great extent, the details of the evolution of galaxies and clusters depend on properties of dark matter. Without knowing those properties, it is difficult to explain how galaxies evolved into the structures observed today. As knowledge of the early universe deepens, I remain optimistic that we will soon know much more about both galaxy formation and dark matter. What we fail to see with our eyes, or detectors, we can occasionally see with our minds, aided by computer graphics. Computers now play a key role in the search for dark matter. Historically, astronomers have focused on observations; now the field has evolved into an experimental science. Today's astronomical experimenters sit neither at lab benches nor at telescopes but at computer terminals. They scrutinize cosmic simulations in which tens of thousands of points, representing stars, gas and dark matter; interact gravitationally over a galaxy's lifetime. A cosmologist can tweak a simulation by adjusting the parameters of dark matter and then watch what happens as virtual galaxies evolve in isolation or in a more realistic, crowded universe.

Old consciousness will not do. Newer one has to be imbibed and cultivated. When people are unaware of their inner wisdom and wholeness, they tend to search for order and security outside themselves. That is my personal experience. There is always search for sanctity of emotion, aesthetic expression, pursuit of harmony, unraveling realms, syncopated ending and stillness of mind and shiver of heart at the mirror broken somewhere are the strains of newer wisdom if it can be called one. Same questions shall be asked and same shall be answer in the similar vein. Uncertain of their place within the hierarchical order of the world, they define themselves based upon their insecurities. Individuals thus become only pieces of their wholeness and like shards of glass from a beautiful vase; they bear little resemblance to their aggregate beauty. Within the hierarchy, many in high positions of power have taken advantage of our collective insecurities in attempting to guide the development of all humankind. They have obscured the direct connection between the individual and the Divine through a variety of means designed to intercede between our inner essence and our divine source. Each individual can come to know them to be free of all forms of institutional and organisational and societal hierarchical control. This is not to imply that we should not trust others or join together in bonds of friendship and community. It is simply a reminder that relative truth is constantly shifting in the hands of those who desire to control. And even when the motive for their controlling behaviors may be of good will, it is still a form of control. When the revealers of "truth" within the hierarchy withhold and suppress information, they are usually positioning themselves to maintain or acquire more power rather than to disseminate empowerment to all. The desire for connection and wholeness is a fuel that drives us to seek out and explore the hierarchical paradigm. This inner longing provides us with the motivation to seek help and guidance from a specific group within the hierarchy, and in so doing, cultivate a sense of belonging and connection. Furthermore, the hierarchical paradigm is a stage whereby we develop a sense of connection to some grand, encompassing vision. This is why the hierarchy nurtures prophets who point toward a greater vision. Spiritual leaders are able to peer deeply beneath the surface reality of life and experience how intricately connected every life form is, and how the composite of all life is intelligent far beyond measure. These visionary leaders can thus interpret reality through their personal abilities to perceive and express life's dimensional depth and limitless intelligence. Yet no one is able to articulate life's full dimensional depth and breadth with the tools of language. They can only, at best, describe their interpretation or their impressions. In actuality, all of us are able to peer beneath the surface reality of life and perceive a unique vision of the universe. We require only time and intention to develop our own interpretations. And this is precisely what many great spiritual leaders have taught. Life's deeper essence is not an absolute to be experienced by the chosen few, but an evolving, dynamic intelligence that wears as many faces as there are life forms. No group or people have the exclusive portal into the universe by which the Divine expresses itself in all its majesty. The portal is open and available to all, because the Divine is within all things. Those recognized as great prophets each produced a vision of the universe beyond what was currently defined by the hierarchy. Because their interpretations were articulated with authority and depth of insight, they became a target of debate among various groups in the hierarchy. This debate then created a polarity of belief. A sympathetic constituency emerged to defend and embellish their leader's interpretation, while established groups held it in contempt of previously held beliefs. Invariably, the leader's vision became confined and shaped into dogma by followers who desired to create a new religion or sect. Thus, this infusion of fresh insight quietly receded into the hands of the hierarchy, where its deeper meaning was obscured by the very fact that it was incorporated into a rigid structure that both protected and promoted it. A new paradigm is emerging that promotes a clear connection of individual consciousness to the compelling features of the Divine without the intervention of a hierarchy. This is when the fables and myths of history step into the light and become known as they were originally intended. This is the time when language will be transformed into a new form of communication that breaks down all barriers of control. Personal transformation, through the awakening of inner wisdom and understanding, is the pathway into wholeness. The transformation paradigm is initiated simply by the recognition that rather than the dependency-inducing ways of the hierarchical paradigm, there are accelerated, independent pathways that bypass the hierarchy and lead to self-mastery. These new pathways lead to the divine wisdom and understanding that is present within all of us. This wisdom can be accessed through the practice of three principles of transformation: seeing the Divine in

all, nurturance of life, and gratitude. The application of these life principles disengages individuals from the controlling elements of the hierarchy, thus initiating the transformation experience. As there are relative truths, there are relative freedoms. As individuals evolve through the hierarchical process, an ever-increasing sense of freedom is gained, yet external forces continue to exert control through limiting language, confining belief systems, and more. These controlling influences lead to anatomical architecture, body, build, cagey carcass, construction, enclosure, fabric, flounce, form, framework, fringe, groundwork, hem, mount, mounting outline, physique, scaffold, scaffolding scheme, setting shell, stage, structure, support, system, trim, trimming, truss- valance continued reliance upon the hierarchy as it unceasingly attempts to impose a sense of inequality between us. The underlying equation of the hierarchical paradigm is: individual + hierarchy = God connection. In the case of personal transformation it is: individual + inner wisdom and understanding divine equality with all. The time has come to integrate the dominant paradigm of the hierarchy with the liberating transformation paradigm. That shall redeem all promises made to oneself. There shall be a Matrix force or balancing force that shall remain UN satisfies until conservativeness is reached. This integration occurs naturally once we have fully explored the two paradigms and develop a synthesis paradigm whereby transformation is attained by searching for connection and wholeness through our own inner wisdom and understanding. It is this combination of release from dependency on the hierarchy and transformation from within that initiates the synthesis paradigm. Once we take responsibility for our own transformation and integration, it does not mean that the hierarchy is to be shunned or avoided. The hierarchy is quite benign as a manipulative force and merely represents one important stage in the journey toward wholeness. What is being set in motion now is the initial preparation for these paradigm shifts. More specifically, these paradigms will be simultaneously played out over the coming years. As always, it will be the choice of each one of us as to which paradigm we embrace in our journey. All of the highest imaginings of humanity are yet unaware of our deepest foundation. We have sought the upper reaches of the building, yet remain unaware of the foundation's design. It is here, at the very core of existence that the Divine is bursting forth with its creative energy while simultaneously reintegrating with its invitation to wholeness. It is here that equality is realized, not in the lofty places of relative truth lodged in the hierarchy, but rather in the deepest part of the foundational plan of life's original source and ultimate destiny. The origin and destiny of existence is the tone of equality in life. Listen for this tone – this frequency of vibration – and follow it back into the very foundation from whence all things arise and return. Consider these words as symbols only. Feel the truth that stands behind these words, and tap into this empowering energy force that reaches out for you. Know it as a tone or vibration – a resonance that waits for you around every corner in which your life will turn. It is a beacon of the Divine gathering itself into the form of language in order to usher you to a place from which you can experience the formless tone of equality – the bypass of limitation. It is the primal language of our divine source that bestows to you the freedom to generate your own deepest beauty in the expression of your highest truth.(Reference: Wikipedia)

## Gopalasamy Model For Procrastinated Conducive, Felicitous Counter Sallies In Two Specified Systems

Abstract

Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of decay rate of perturbations. Sole intention, main objective and primary concern of

the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are counted in the same vein as that of predator prey. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 11. Preliminary remarks

Almost all the equations of population ecology have been derived starting from the following simple format (Note this can be used for the development of the creed and type of systemic types in case of cosmology and astrophysics and fiancé and econometrics. To give an example of finance, we can think of a account of the company giving rise to various accounts of its employees, and that is the contribution of the single account for the development of the total system types of the given configuration). It may be account, particle, atom, molecule, proposition, problem or any such thing. We shall delineate and disseminate information and particulars as and when the model shall be applied to the problem in question.

$$\frac{dN(t)}{dt} = \text{an individual's contribution}$$

*to population growth* $N(t)$

(1.1)

where $N(t)$ denotes the size of a population of a single species (or biomass) with overlapping generations at time $t$, and $dN(t)/dt$ is the rate of change of the population size at time $t$. Subsequently one makes an assumption regarding the factor in the brackets in (1.1); in particular if one assumes that an individual entity's contribution to the population growth is denoted by a function say, f(t, N), defined suitably for all t>0, N then one obtains from (1.1) the so called Kolmogorov formulation in the form

$$\frac{dN(t)}{dt} = f(t, N(t))\, N(t).$$

(1.2)

Various choices of $f$ together with some ecologically plausible assumptions such as the temporal constancy of the environment and density dependent effects in $f$ lead to various equations. For instance if $f(t, N) \equiv r$ (a positive constant) one obtains the Malthusian formulation in the form

$$\frac{dN(t)}{dt} = rN(t)$$

And if one assumes $f(t, N) \equiv r - (r/K) N$ for some positive constants r and K one gets the familiar logistic equation of population biology

$$\frac{dN(t)}{dt} = rN(t) \left\{ 1 - \frac{N(t)}{K} \right\}$$

In order to device models with time delayed effects, such effects are introduced in the average growth rate $F(t, N)$ of (1.2); for instance if $f(t, N) \equiv r - (r/K) N(t - \tau)$ for some constant $\tau > 0$ we obtain Hutchinson's[19]delay loistic equation

$$\left\{ \frac{dN(t)}{dt} = rN(t) \left\{ 1 - \frac{N(t-\tau)}{K} \right\} \right. .$$

(1.3)

If for instance

$$F(t, N) = r - (r/K) N(t) - \int_{-\infty}^{t} k(t - s) F(N(s))\, ds$$

Where $k$ is a piecewise continuous nonnegative valued function such that $\int_0^\infty k(s)\,ds = 1$ with some appropriate function $F$, we will get Volterra's model of a population which pollutes its environment and pollution has an accumulative and toxic effect; such an equation is of the form

$$\frac{dN(t)}{dt} = N(t) = \left\{ r - (r/K) N(t) - \int_{-\infty}^{t} k(t-s) F(N(s))\,ds \right\} \quad (1.4)$$

Which has been studied in different formats by several authors (Cushing [12], Cohen *et al.* [10], and Gopalasamy [16])

One of the significant advantages of deriving various models from the prototype (1.2) is that if the initial; conditions (initial population sizes) are nonnegative then the nonnegativity of the population size $N(t)$ for $t > 0$ follows from the fact that solutions of (1.2) satisfy

$$N(t) = N(0) \exp\left[\int_0^t f(s, N(s))ds\right]$$

Which implies that $N(t) \geq 0$ if $N(0) \geq 0$.

While the time delayed logistic equation (1.3) explains some of the observed oscillations of certain populations, its derivation has not been entirely satisfactory. For instance it has been pointed out by Ricklefs ([27], page 488) that "time lags might be expected to occur primarily in stabilizing processes involving production rather than death because death is an immediate response to environmental change". Recently Cushing [13] has examined the effects of response and maturation delays in a number of single and multi system models; as it has been noted by Cushing [13], the models involving delayed reproduction or not necessarily derivable from the Kolmogorov type systems of the form (1.2) by suitably defining the average growth rate $(1/N)(dN/dt) = f(t, N(t))$. While Cushing [13] has been mostly concerned with the possibility of delay induced oscillations we will consider a class of delays in reproduction and interspecific interaction which maintain the stability of the systems.

Our starting point to drive models with delayed reproduction is the following balance equation which assumes that there is no immigration or emigration. Account, particle, atom, molecule, proposition, problem et al

$$\frac{dN(t)}{dt} = \text{birth rate} - \text{death rate}. \quad (1.5)$$

Now for instance if we consider a population of adult flies then the birth rate at time $t$ depends on the adult population at time $t^{-}\ \tau$ where $\tau$ is the time required for the larvae to become adult flies. If the birth and death rates are governed by density dependent factors then we have from (1.5) the following time delayed model

$$\frac{dN(t)}{dt} = b\big(N(t^{-}\tau)\big) - m(N(t))$$

(1.6)

Where the functions $b$ and $m$ denote the density dependent birth and death rates of Account, particle, atom, molecule, proposition, problem et al respectively. If the time delay is continuously distributed then one can consider in the place of (1.6) an equation of the form

$$\frac{dN(t)}{dt} = b\int_{-\infty}^{t} k\{(t-s)N(s)ds - m(N(t))$$

Or equivalently

$$\left\{\frac{dN(t)}{dt} = b\int_{0}^{\infty} k(s)\right\}N(t-s)ds - m(N(t))$$

(1.7)

Where the delay kernel $k$ denotes the weight of the delayed of the delayed effects.

It is not always obvious that suitable nonnegative initial conditions for (1.6) and (1.7) will imply the nonnegativity of the solutions of (1.6) and (1.7). The purpose of this article is to show that in two species systems of competitive and prey- predator communities, delayed reproduction and delayed interactions donot destabilise the systems if the self- regulating negative feedback effects are sufficiently strong compared to interspecific interactive effects and if the self regulating effects are realised with no time delays. When we say that time delays do not destabilise, we will mean that delays do not render an asymptotically (locally) stable steady state unstable in such a manner that the loss of stability either leads

To a delay – induced bifurcation to persistent and undamped oscillations (as in the case of Hopf- bifurcation) or the relevant linear variational system has unbounded solutions; that is a mathematical instability is not induced by delays in the sense, that the roots of the characteristic equation are pushed to the imaginary axis or to the right of it on the complex plane. However delays can make an otherwise stable system less stable which means that the rate of convergence to zero of solutions of the variational system can decrease with an increase of delay and the variational system can even oscillate (perhaps violently) before converging to the trivial state. Such a reduction in stability can

happen even when all the roots of the characteristic equation have negative real parts. For our purposes such a reduced stability is still stability. In order to see the effects of delays on such less stable systems we will propose below a method of estimating the rate of decay of local perturbations based on the variational system and indirectly relying on the characteristic equation. For a special class of two species models with a single delay we have obtained (Gopalsamy [16]) estimates of decay rates (or return times) usingthe characteristic equation; such a procedure is almost impossible in the case of many species systems with several delays. The method we propose in Section 6 below for the estimate of decay rates can be adapted to *n*- species systems with $n^2$ discrete delays.

Let $x_1(t)$ and $x_2(t)$ denote the population sizes (or biomasses) of two systems competing for a common pool of resources.. Let $b_i$ and $m_i$ $(i = 1, 2)$ denote the respective density dependent birth (potential) and death rates so that in the absence of any time delays the population sizes are governed by (of Account, particle, atom, molecule, proposition, problem et al). In addition as usual, we assume that the system studied follow Haimovici's equations, for which the dove tailing explanation has been appended

$$\left.\begin{aligned}\frac{dx_1(t)}{dt} &= b_1\left(x_1(t)\right) - m_1(x_1(t)), x_2(t))\\ \frac{dx_2(t)}{dt} &= b_2\left(x_2(t)\right) - m_2(x_1(t)), x_2(t))\end{aligned}\right\} \quad (2.1)$$

In order to make the system (2.1) represent a competition system we will make the following assumptions on the birth and death rates; $b_1$, $m_1$ $(i = 1, 2)$ are continuous with continuous partial derivatives for all $x_1 \geq 0$ $(I = 1, 2)$; also we assume (see for instance Rescigno and Richardson [26]) $\frac{\partial b_i}{\partial x_1} > 0, \frac{\partial m_1}{\partial x_j} > 0$ for $x_i > 0$ $(i, j = 1, 2)$, (2.2)

(xvi) $b_1(0) = 0$, $m_1(0, x_2) \equiv 0, b_2(0) = 0$, $m_2(x_1, 0) \equiv 0$, (2.3)For some $x_{1,}^* > 0, x_{2,}^* > 0$ we have

$$\left.\begin{aligned}b_1(x_1^*, -m_1(x_1^*, 0) &= 0,\\ b_2(x_2^*, -m_1(0, x_2^*) &= 0,\end{aligned}\right\} \quad (2.4)$$

(xvii) There exist positive constants $\delta_1$, $\delta_2$ such that

$b_1(\delta_1) - m_1(\delta_1, x_2) < 0$,

$b_2(\delta_2) - m_2(x_1, \delta_2) < 0$, for $x_1 \geq 0$, $x_2 \geq 0$,

(2.5)

(xviii) There exists a point $(\alpha, \beta), \alpha > 0, \beta > 0$ such that

$$b_1(\alpha) - m_1(\alpha, \beta) = 0 = b_2(\beta) - m_2(\alpha, \beta).$$

(2.6)

The conditions on $b_i$ in (2.2) mean that the potential (not average) birth rates are density dependent and any negative density dependent feedbacks to the opened or birth rate of Account, particle, atom, molecule, proposition, problem et al rates are included in the destruction rates; the conditions on $m_i$ indicate intraspecific and interspecific confederation needs and necessities. (2.3) implies that (0, 0) is a (trivial) steady state of the system (2.1) as it is in the case of all ecological models; (2.4) means that in the absence of either of the two species the other one has a nontrivial steady state; (2.5) implies that each of the two populations cannot grow unboundedly since when $x_1(t) > \delta_i, \frac{dx_1}{dt} < 0$. (2.6) means that there exists a nontrivial steady state $(\alpha, \beta)$ in the interior of the positive quadrant of the $(x_1, x_2)$- plane. These assumptions are standard and are minimal in the opinion of the author.

An immediate example of (2.1) –(2.6) is the familiar Volterra –Lotka competition system described by

$$\frac{dx_1(t)}{dt} = r_1 x_1(t) - a_{11} x_1^2(t) - a_{12} x_1(t) x_2(t),$$

$$\frac{dx_2(t)}{dt} = r_2 x_2(t) - a_{21} x_1(t) x_2(t) - a_{22} x_2^2(t),$$

(2.7)

Where $r, a_{ij}, (i, j = 1,2)$ are positive constants satisfying the condition

$$\frac{a_{11}}{a_{21}} > \frac{r_1}{r_2} > \frac{a_{12}}{a_{22}}.$$

One of the first questions to be resolved for (2.1) is the following: if $x_i(0) > 0$, $x_2(0) > 0$ does it follow that the corresponding solutions of (2.1) satisfy the condition $x_1(t) \geq 0$, $x_2(t) \geq 0$ for those $t$ for which a solution of (2.1) exists? In the $(x_1, x_2)$ phase plane of (2.1) we have (0, 0), $(x_1^*, 0)$, $(0, x_2^*)$, $(\alpha, \beta)$ as the steady states of (2.1) and also it will follow from (2.3) that the co ordinate axes

$x_1 = 0$ and $x_2 = 0$ are invariant sets f (2.1) ; such invariance of the coordinate axes will imply that if $x_1(0) > 0$, $x_2(0) > 0$ then the corresponding population trajectory $(x_1(t), x_2(t))$ of (2.1) cannot reach the outside of the nonnegative quadrant $x_1 \geq 0$, $x_2 \geq 0$.

Our plan is to derive first a set of sufficient conditions for the coexistence of the two competing species and then in the next section show that under the same conditions the two species can coexist if there are time delays in the reproduction and interspecific interactions. We recall that the two competing species described by (2.1) are said to coexist if and only if the nontrivial steady state $(\alpha, \beta)$ is (at least locally) asymptotically stable. To examine such a stability of $(\alpha, \beta)$ we derive the corresponding linear variational system for the perturbations $X_1, X_2$ where

$$x_1(t) \equiv \alpha + X_1(t), \qquad x_2(t) \equiv \beta + X_2(t)$$

Such a variational system is found to be

$$\left.\begin{aligned} \frac{dX_1(t)}{dt} &= \beta_{11} X_1(t) + \beta_{12} X_2(t), \\ \frac{dX_2(t)}{dt} &= \beta_{21} X_1(t) + \beta_{22} X_2(t), \end{aligned}\right\} \qquad (2.8)$$

Where

$$\beta_{1J} = \left(\frac{\partial b_1}{\partial x_J} - \frac{\partial m_1}{\partial x_j}\right) \text{ evaluated at } (\alpha, \beta) \qquad (i, j = 1,2).$$

(2.9)

The steady state $(\alpha, \beta)$ of (2.1) is locally asymptotically stable if and only if the trivial solution (0, 0) of (2.8) is asymptotically stable. We can derive the following result from elementary considerations.

In the system (2.1)-(2.6) assume the following:

$$\left.\begin{aligned} \frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1}, \\ \frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1}, \end{aligned}\right\} \qquad \text{at } (\alpha, \beta)$$

(2.10)

Then the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (2.1) – (2.6).

The proof follows easily if we note that the characteristic equation associated with the linear system (2.8) is given by

Det $\begin{pmatrix} \lambda-\beta_{11} & -\beta_{12} \\ -\beta_{21} & \lambda-\beta_{22} \end{pmatrix} = 0$

Or equivalently,

$$\lambda^2 - \lambda\,(\beta_{11} + \beta_{22}) + \beta_{11}\beta_{22} - \beta_{12}\beta_{21} = 0. \qquad (2.11)$$

It follows from (2.2) and (2.10) that,

$$\beta_{11} < 0, \beta_{11}\beta_{22} - \beta_{12}\beta_{21} > 0$$

And hence the roots of (2.11) have negative real parts implying the asymptotic stability of the trivial solution of (2.8)

If apply the condition (2.10) to the Volterra-Lotka model (2.7) we note that when a nontrivial steady state say $(\alpha, \beta)$ exists for (2.7), (2.10) leads to

$$a_{11}\alpha > a_{21}\,\beta \text{ and } a_{22}\,\beta > a_{12}\,\alpha,$$

Which together imply $a_{11}\,a_{22} > a_{12}a_{21}$. If $(\alpha, \beta)$ exists it is known that $a_{11}\,a_{22} > a_{12}a_{21} > 0$ is a sufficient condition for the (local) asymptotic stability of $(\alpha, \beta)$.

The condition (2.10) which guarantees the coexistence of competing populations in (2.1)-(2.6) simply means that the intraspecific self – regulating feedback effects $\partial m_i / \partial x_1\ (i = 1, 2)$ bare higher than its own positive feedback $\partial b_1 / \partial x_1\ (i = 1, 2)$ together with its negative (competitive) effects on its competitor $\partial m_I / \partial x_J\ (i,\ j = 1, 2; i \neq j)$.

## 12. Coexistence under delayed production and interactions, transactions

We will now establish that two competing species with a sufficiently strong intraspecific regulation can coexist in the presence of delayed reproduction and delayed interspecific interactions. It is well known that if there are delays in the self – regulating intraspecific effects, stationary coexistence may not always be possible (Gopalsamy [16], Gopalsamy and aggarwala [15], Shibata and Saito [28]). Delayed reproduction is more general and widespread in most biological populations. To build a competition model with such delays we proceed as follows.

Let $\tau_{ij}\ (i, j = 1, 2)$ be a set of nonnegative constants with $\tau = \max \tau_{ij}\ , i, j = 1, 2,$ and suppose that the two ompeting species display delayed reproduction and delayed interspecific interactions with no delays in the self – regulating negative feedback effects. Such a competition system in a constant environment can be modeled by the autonomous delay differential system

$$\left.\begin{aligned} \frac{dx_1(t)}{dt} &= b_1\big(x_1(t-\tau_{11})\big) - m_1\big(x_1(t), x_2(t-\tau_{12})\big), \\ \frac{dx_2(t)}{dt} &= b_2\big(x_2(t-\tau_{22})\big) - m\big(x_1(t-\tau_{21}), x_2(t)\big), \end{aligned}\right\} \qquad (3.1)$$

Where both the creation and destruction rates and the morality rates $b_1$ and $m_1\ (i = 1, 2)$ satisfy the same conditions (2.2)-(2.6) as in the previous section. Along with (3.1) we will

Suppose that the initial population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) sizes are provided in the form

$$x_i(s) = \phi_i(s) \geq 0, \qquad s \in [-\tau, 0], \qquad i = 1,2. \tag{3.2}$$

(Note that the specification of $\phi_i$ need not be on the entire interval$[-\tau, 0]$; but such a hypothesis is convenient.) Since (3.1) is not of the standard Kolmogorov-type, we have to establish that the solutions of (3.1) will remain nonnegative for nonnegative initial conditions.

To prove the nonnegativity of solutions of (3.1)-(3.2), let $\left\{\psi_1^{(n)}(s)\right\}, s \in [-\tau, 0], i = 1,2; (n = 1,2, \ldots)$ be a sequence of strictly positive continuous functions such that

$$\lim_{n\to\infty} \psi_1^{(n)}(s), = \psi_1(s), \qquad s \in [-\tau, 0], \qquad i = 1,2 \tag{3.3}$$

Let $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ be the solution of (3.1) corresponding to the initial condition

$$\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \left\{\psi_1^{(n)}(s), \psi_2^{(n)}(s)\right\}, \quad s \in [-\tau, 0],. \tag{3.4}$$

Consider the solution $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ for $t \in [0, \tau]$ (assuming it exists on the whole of $[0, \tau]$). Suppose $x_1^{(n)}$ does not remain nonnegative for all $[0, \tau^*]$; then there exists a $t^* = [0, \tau^*](\tau^* =$ the positive minimum of $\tau_{ij}\ (i, j = 1, 2))$ for which

$$x_1^{(n)}(t^*) = 0 \ and \ x_1^{(n)}(t)(t^* - \tau_{11})) > 0.$$

It will follow from (3.1) and the positivity of the initial condition that

$$\frac{dx_1^{(n)}(t^*)}{dt} = b_2\left(x_1^{(n)}(t^* - \tau_{11})\right) > 0 \tag{3.5}$$

Where we have used the properties of $b_i$ and $m_i$ from (2.2)-(2.6); it is found that (3.5) contradicts the definition of $t^*$. Thus $x_1^{(n)}(t) > 0$ on $[0, \tau^*]$; similarly $x_2^{(n)}(t) > 0\ on\ [0, \tau^*]$. We can repeat this procedure for intervals of the form $[\tau^*, 2\tau^*], [2\tau^*, 3\tau^*]\ etc$. Thus it will follow that so long as $(x_1^{(n)}(t), x_2^{(n)}(t))$ exists, we have

$$x_1^{(n)}(t) \geq 0\ x_2^{(n)}(t) \geq 0.$$

If we now consider the limits as $n \to \infty$, we get

$$\lim_{n\to\infty}\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \{x_1(t), x_2(t)\} \tag{3.2}$$

With $x_1(t) \geq 0, , x_2(t) \geq 0$ where $(x_1(t), x_2(t))$ is the solution of (3.1)-(3.2) and this is a consequence of the continuous dependence of solutions of (3.1) on initial conditions (Hale [17], page 41).

A second question for (3.1) with (2.2)-(2.6) is whether solutions of (3.1) exist for all $t > 0$? suppose now that solutions of (3.1) do not exist for all $t$ ; then there exists a $t_1 > 0$ such that at least for one of either of $x_1$ *or* $x_2$ we have

$$\lim_{t \to t_1 -} x_i\ (t) = \infty. \tag{3.7}$$

To be specific let us suppose that $\lim_{t \to t_1 -} x_i\ (t) = \infty$; let $t_2$ be the first instant of time for which $x_1(t_2) = \delta_1, t_2 < t_1$. It will follow from (3.1) that

$$\frac{dx_1(t_2)}{dt} = b_1\big(x_1\ (t_2 - \tau_{11})\big) - m_1\big(x_1(t_2), x_2\ (t_2 - \tau_{12})\big)$$

$$< b_1\big(x_1\ (t_2)\big) - m_1\big(x_1(t_2), x_2\ (t_2 - \tau_{12})\big)$$

$$< b_1(\delta_1) - m_1\big(\delta_1, x\ (t_2 - \tau_{12})\big) < 0 \tag{3.8}$$

Which contradicts the definition of $t_2$. Thus solutions of (3.1), (2.2) – (2.6) exist for all $t$ .

We are now ready to consider the effects of time delays in reproduction and interspecific interactions on the coexistence of the two competing species. The steady state $(\alpha, \beta)$ of (2.1) is again a steady state of (3.1). To examine the (local) asymptotic stability of $(\alpha, \beta)$ in (3.1) we let

$$x_1(t) = \alpha + Y_1(t), \quad x_2(t) = \beta + Y_2(t)$$

In (3.1) and derive that the linear variational system governing the perturbations $Y_1, Y_2$ is given by

$$\frac{dY_1\ (t)}{dt} = -m_{11}Y_1(t) + b_{11}Y_1\ (t - \tau_{11}) - m_{12}Y_2\ (t - \tau_{12}),$$

$$\frac{dY_2\ (t)}{dt} = -m_{22}Y_2(t) + b_{21}Y_1\ (t - \tau_{21}) - m_{22}Y_2\ (t - \tau_{22}),$$

Where

$$m_{ij} = \frac{\partial m_1}{\partial x_j}, \quad b_{ii} = \frac{\partial b_i}{\partial x_i}\ at\ (\alpha, \beta). \tag{3.10}$$

It is well known (Bellman and Cooke [3], page 336) that the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) if the trivial solution of the variational system (3.9) is asymptotically stable. We formulate our first result in the following form whose proof is given in the appendix.

*THEOREM 3.1. Assume that in the two systemic types competition model system (3.1) with delayed density dependent birth and death rates $b_i$ and $m_i$ $(i = 1, 2)$ satisfy the conditions (2.2)-(2.6). Suppose that the conditions (2.10) hold. Then the nontrivial steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) whatever the delays be $\tau_{ij}$ $(i = 1, 2)$.*

Let $y_1(t) and\ y_2(t)$ respectively denote the population sizes (or biomasses) of a prey species and a predator species which feeds exclusively on the single prey species in a temporally constant environment. In the format of (1.5) the following system will describe the rates of change of the two population sizes:

$$\frac{dy_1(t)}{dt} = f_1(y_1(t)) - g_1(y_1(t), y_2(t)),$$

$$\frac{dy_2(t)}{dt} = f_2(y_1(t)) - y_2(t) - g_2(y_2(t)), \quad (4.1)$$

Where $f_i\ and\ g_i\ i = 1,2)$ respectively denote the birth and death rates of the two species; in order to make (4.1) denote the dynamics of a prey – predator species we will assume that $f_i\ and\ g_i$ satisfy the following conditions (see for instance Rescigno and Richardson [26]); $f_i\ and\ g_i$ are continuous with continuous partial derivatives for all $y_i \geq 0\ (i = 1, 2); also\ we\ require$

$$\frac{\partial f_1}{\partial y_1} > 0; \quad \frac{\partial g_1}{\partial y_1} \geq 0; \quad \frac{\partial g_1}{\partial y_2} \geq 0,$$

$$\frac{\partial f_2}{\partial y_1} \geq 0; \quad \frac{\partial g_2}{\partial y_2} \geq 0; \quad \frac{\partial g_2}{\partial y_2} \geq 0, \quad \text{Where } y_1 \geq 0, y_1 \geq 0, \quad (4.2)$$

(xxi)

$$f_1(0) = 0, \quad g_1(0, y_2 \equiv 0,$$

$$f_2(0, y_2) = 0 = f_2(y_1, 0); \quad g_2(0) = 0, \quad (4.3)$$

(xxii) For some constant $y_1^* > 0$ we have

$$f_1(y_1^*) - g_1(y_1^* 0) = 0, \quad (4.4)$$

(xxiii) There exists a point $(\alpha^*\beta^*), \alpha^* > 0, \beta^* > 0$ such that

$$f_1(\alpha^*) - g_1(\alpha^*, \beta^*) = 0,$$

$$f_2(\alpha^*, \beta^*) - g_2(\beta^*) = 0, \quad (4.5)$$

(xxiv) There exist positive constants $n_1, n_2$ such that

$$f_1(\eta_1) - g_1(\eta_1, y_2) < 0, y_2 \geq 0,$$

$$f_2(\eta_1, \eta_2) - g_2(\eta_2) < 0 \quad (4.6)$$

The conditions (4.2) mean that the system (4.1) is of the oppugnacy systems type; (4.3) will imply that (0,0) is a trivial steady state of (4.1); (4.4) means that in the absence of the predator, the prey species has a nontrivial (positive) steady state;

(4.5) implies that (4.1) has a nontrivial steady state in the interior of the positive quadrant of the state species of (4.1); (4.6) can be shown (see below) to imply that no credit type and debit type system types can grow unboundedly.

An example of (4.1)-(4.6) is the familiar Volterra-Lotka prey-predator system

$$\frac{dy_l(t)}{dt} = r_l y_l(t) - [a_{ll} y_l^2(t) + a_{12} y_l(t) y_2(t)],$$

$$\frac{dy_2(t)}{dt} = a_{12} y_1(t) y_2(t) - [r_2 y_2(t) + a_{22} y_2^2(t)],$$

Where $r_{l,}\ a_{ij}\ (i,j=1,2)$are positive costants such that

$$a_{21}r_1 - r_2a_{11} > 0.$$

As in the previous case of competition we have to verify that when $y_1(0) > 0, y_2(0) > 0$ the corresponding solution of (4.1) will be nonnegative for $t$ 0 for which the solution exist. Since the co-ordinate axes $y_1 = 0, y_2 = 0$ are invariant sets for (401) by virtue of (4.3), it will follow that the solution of (4.1) starting from the interior of the positive quadrant cannot reach the outside of that quadrant. This is a consequence of the invariance of the co-ordinate axes for (4.1)

To prove that solutions of (4.1)-(4.6) exist for all t 0 let us suppose that $y_l 0\ and\ y_2(0) satisfy\ 0 < y_l\ (0)$ satisfy $0< y_l(0) < \eta_l, 0 < y_2(0) < \eta_2$ ( on the other hand if either one of $y_l(0)$ $\eta_l and\ y_2(0)$ $\eta_l$ or both hold, it will follow that the corresponding $y_l(t)$ will decrese by virtue of (4.6) at least for small $t$ 0, ). Suppose $y_l(t)$ does not exist for all $t$ 0; then for some $t_l < \infty\ lim_{t\to t_l-} y_l(t) = \infty$and let $t_2 > 0$ be the first instant for which $y_1(t_2) = \eta_1$; then we have from (4.1) and (4.6) that

$$\frac{dy_1(t_2)}{dt} = f_1(\eta_1) - g_1(\eta_{1,}\ y_2) < 0\ (assumimg\ y_1(0) < \eta_1)$$

Which contradicts the definition of $t_2$ showing that $y_1(t)$ exists for all $t$ 0; if $y_2(t)$ is not defined for all $t$ 0, then for some $t_3$ $\infty$ we will have $lim_{t-t_3-} y_2(t) = \infty$;now let $t_4$ be the first instant for which $y_2(t_4) = \eta_2$; then from the properties of $f_2$ we have

$$\frac{dy_2(t_4)}{dt} = f_2(y_1(t_4), \eta_2) - g_2(\eta_2)\ f_2(\eta_1, \eta_2) - g_2(\eta_2)$$

$$< 0 (by\ (4.6) assuming\ 0 < y_1(0) < \eta_1, 0 < y_2(0) < \eta_2)$$

Which again contradicts the definition of $t_4$. Thus $y_2(t)$ is defined for all $t$ 0.

We will now derive a set of sufficient condition for the (local) asymptotic stability of the steady state $(\alpha^*, \beta^*)$ of (4.1). in the next section we will show that under the same conditions and in the presence of various delays , the steady state $(\alpha^*, \beta^*)$remains (locally) asymptotically stable.

To examine the asymptotic stability of the nontrivial steady state $(\alpha^*, \beta^*)$ in (4.1) we let

$$y_1(t) = \alpha^* + U_1(t), \qquad y_2(t) = \beta^* + U_2(t)$$

In (4.1) and derive the liner variational system in the perturbations $U_1, U_2$ as follows:

$$\left.\begin{aligned} \frac{dU_1(t)}{dt} &= \frac{\partial f_1}{\partial y_1}U_1(t) - \frac{\partial g_1}{\partial y_1}U_1(t) - \frac{\partial g_1}{\partial y_2}U_2(t), \\ \frac{dU_2(t)}{dt} &= \frac{\partial f_2}{\partial y_1}U_1(t) - \frac{\partial f_2}{\partial y_2}U_2(t) - \frac{\partial g_2}{\partial y_2}U_2(t), \end{aligned}\right\} \quad (4.7)$$

Where all the partial derivatives in (4.7) are evaluated at $(\alpha^*, \beta^*)$. The steady state $(\alpha^*, \beta^*)$ of (4.1) is (locally) asymptotically stable if and only if the trivial solution of (4.7) is asymptotically stable which will be the case if and only if the roots of the characteristic equation

$$\text{Det} \begin{pmatrix} \lambda - (f_{11} - g_{11}) & g_{12} \\ -f_{21} & \lambda - (f_{22} - g_{22}) \end{pmatrix} = 0,$$

Where $f_{ij} = \frac{\partial f_1}{\partial y_l}$, $g_{ij} = \frac{\partial g_1}{\partial y_j}$, at $(\alpha^*, \beta^*)$ (4.8)

$(i,j = 1, 2)$

Have negative real parts. If $\lambda_1, \lambda_2$ are the riits of (4.8) then

$$\lambda_1 + \lambda_2 = (f_{11} - g_{11}) + (f_{22} - g_{22}), \quad (4.9)$$

$$\lambda_1 \lambda_2 = (f_{11} - g_{11})(f_{22} - g_{22}) + f_{21} g_{12}.$$

The proof of the following result is an immediate consequence of (4.7)-(4.9)

THEOREM 4.1 *In the CREDIT-DEBIT system (4.1)-(4.6) assumes that the partial derivatives satisfy the following additional conditions:*

$$\frac{\partial g_1}{\partial y_1} > \frac{\partial f_1}{\partial y_1} + \frac{\partial f_2}{\partial y_1}, \quad \frac{\partial g_2}{\partial y_2} > \frac{\partial f_2}{\partial y_2} + \frac{\partial g_1}{\partial y_2}, \quad \text{Evaluated at } (\alpha^*, \beta^*), \quad (4.10)$$

*Then the steady state $(\alpha^*, \beta^*)$, is (locally) asymptotically stable for the prey- predator system.*

## 10. Delays in reproduction, predation and gestation

A number of authors have proposed and analysed several models of prey- predator systems with time delays. Most of the models have been derived from the so called models of the Kolmogorov-type f (1.2) with time delays incorporated in the average growth rates of the prey and predator species. Volterra [32] proposed the following model for a prey-predator system. The equations are used for the formulation of Credit-debit system of atoms, accounts, economic variable, cosmological variable, astrophysical; variables, and qubit systems for which the model is being applied.

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \delta_1 N_1(t) - \int_{\infty}^{t} F_1(t-s) N_2(s) ds],$$

(5.1)

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 + \delta_2 N_1(t) - \int_{\infty}^{t} F_2(t-s)N_1(s)ds],$$

Where $r_1, r_2, \delta_1, \delta_2$ *are positive constant and* $F_1, F_2$ is nonnegative continuous delay kernel defined and integrable on $[0, \infty]$. (5.1) has no predator self-regulating negative feedback. Brelot [6] considered a modified format of (5.1) in the form

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \lambda_1 N_1(t) - \delta_1 N_2(t) - \int_{-\infty}^{t} F_1(t-s)N_2(s)ds],$$

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 - \lambda_2 N_2(t) - \delta_2 N_1(t) - \int_{-\infty}^{t} F_2(t-s)N_1(s)ds], \quad (5.2)$$

Which has term representing intraspecific competition among both predator and prey species. Under suitable condition on the various parameters in (5.2) one can show that (5.2) has a (locally) asymptotically stable steady state. A number of intergrodifferential equation models of the types (5.1) and (5.2) have been investigated by Cushing [12]

Starting from Hutchinson's (1948)[19] delay logistic equation, May[23] has proposed the following system

$$\frac{dN_1(t)}{dt} = rN_1(t)\left[1 - \frac{N_1(t-\tau)}{K}\right] - \alpha N_1(t)N_2(t),$$

$$\frac{dN_2(t)}{dt} = -bN_2(t) + \beta N_2(t)N_1(t), \quad (5.3)$$

Where r, $\tau$, K, $\alpha, \beta, b$ are positive constants; (5.3) contains a single discrete delay; one can modify (5.3) and incorporate a continuously distributed delay in it so that (5.3) can be written as (May [23])

$$\frac{dN_1(t)}{dt} = rN_1(t)[1 - \int_{-\infty}^{t} F(t-s)N_1(s)ds] - \alpha N_1(t)N_2(t),$$

$$\frac{dN_2(t)}{dt} = -bN_2(t) + \beta N_1(t)N_2(t). \quad (5.4)$$

The first model of a prey-predator type which departs from the Kolmogorov-type formulation is due to Wangersky and Cunningham [33] who considered the model system

$$\frac{dN_1(t)}{dt} = a_1N_1(t) - b_1N_1^2(t) - c_1N_1(t) - N_2(t),$$

$$\frac{dN_2(t)}{dt} = a_2N_2(t) + c_2N_1(t-\tau)N_2(t-\tau), \quad (5.5)$$

Where $a_1, a_2, b_1, c_1, c_2, \tau$ are positive constants; (5.5) means that a duration of $\tau$ time units elapses when an individual prey is killed and the moment when the corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) Recently Cushing [13] has considered a model of the form

$$\frac{dN_1(t)}{dt} = rN_1(t)\left[1 - \frac{N_1(t)}{k} - \alpha N_2(t)\right],$$

$$\frac{dN_2(t)}{dt} = -d_2^*(t) + bN_1(t)\int_0^\infty \delta(a)N_2(t-a)e^{-d*a}\,da,$$

(5.6)

Where $r, k, \alpha, d^*, b$ are positive constants and $\delta$ is related to the age dependent fecundity of the predator species.

With the models (5.1)-(5.6) in the background let us now consider a time delayed model similar to (4.1) in the following form:

$$\frac{dN_1(t)}{dt} = f_1\big(N_1(t-\tau_{11})\big) - g_1\big(N_1(t), N_2(t-\tau_{12})\big),$$

$$\frac{dN_2(t)}{dt} = f_2\big(N_1(t-\tau_{21})\big), N_2(t-\tau_{12}) - g_1(N_2(t)),$$

(5.7)

Where $\tau_{ij}$ $(i,j=1,2)$ are nonnegative constants with $\tau = \max \tau_{ij}, i,j = 1,2$; the birth and death rate functions $f_1, f_2, g_1, g_2$ satisfy the same conditions as in (4.1)-(4.6). Along with (5.7) we specify the initial conditions as follows:

$$N_1(s) = \phi_1(s) > 0,$$

$$s \in [-\tau, 0].$$

$$N_1(s) = \phi_1(s) > 0,$$

We3 will assume that in addition to (4.2) we have for (5.7) that

$\frac{\partial f_2}{\partial N_1} > 0, \frac{\partial f_2}{\partial N_2} > 0, f_2(N_2, N_2) > 0$ *whenever* $N_1 > 0,\ N_2 > 0.$ (5.9)

It is not difficult to show by the methods previously used for the competition system that whenever $N_1\ (s) > 0$ and $N_2\ (s) > 0$ for $s \in [-\tau, 0\ ]$, solution of (5.7) exist for all $t \geq 0$ and remain nonnegative for all $t \geq 0$. also $(\alpha^*, \beta^*)$ of (4.5) is a steady state of (5.7).

We will discuss the form and meaning of the delays in (5.7) in the next section. The effects of the various delays in (5.7) on the (local) asymptotic stability of $(\alpha^*, \beta^*)$ can be derived from the following theorem whose proof is quite similar to that of Theorem 3.1 and hence the details of the proof are omitted.

THEOREM 5.1 *In the time delayed prey-predator system(* corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) *) (5.7) let the conditions (4.2)-(4.6), (5.8) and (5.9) hold for the birth and death rates; furthermore suppose the condition (4.10) hold. Then for all nonnegative delays in (5.7), its nontrivial steady state $(\alpha^*, \beta^*)$ is (locally) asymptotically stable.*

## 6. Estimation of decay rates

We will first derive a convention sufficient condition for the asymptotic stability of the trivial solution of a linear system of the form

$$\left.\begin{aligned} \frac{dY_1(t)}{dt} &= -m_{11}\, Y_1(t) + b_{11} Y_1(t-\tau_{11}) - m_{12} Y_2(t-\tau_{12}), \\ \frac{dY_2(t)}{dt} &= -m_{22}\, Y_1(t) + m_{21} Y_1(t-\tau_{21}) - b_{22} Y_2(t-\tau_{22}), \end{aligned}\right\} \quad t > 0, \quad (6.1)$$

Where $m_{1J}, b_{i1}\ (i, j = 1,2)$ are positive constants while $\tau_{1J}(i, j = 1,2)$ are non negative with $\tau = \max(\tau_{1J})(i, j = 1,2$.as is shown in the appendix, the following system dominates (6.1*)*

$$\left.\begin{aligned} \frac{dZ_1(t)}{dt} &= -m_{11}\, Z_1(t) + b_{11} Z_1(t-\tau_{11}) - m_{12} Z_2(t-\tau_{12}), \\ \frac{dZ_2(t)}{dt} &= -m_{22}\, Z_1(t) + m_{21} Z_1(t-\tau_{21}) - b_{22} Z_2(t-\tau_{22}), \end{aligned}\right\} \quad t > 0, \quad (6.2)$$

$Z_i(s) = |Y_i(s)|, \qquad i = 1,2\ \ for\ s \leq 0\ \ for\ which\ Y_i(s)\ is\ defined.$

Thus the asymptotic stability of the trivial solution of (6.2) implies that of (6.1). it is also shown in the appendix that the solutions of (6.2) are nonnegative for all t $\geq 0$. we will use the following vector-matrix notation:

$$Z = (Z_1, Z_2)^{\tau}, \qquad ||Z|| = |Z_1(t)| + |Z_2(t)| = Z_1(t) + Z_2(t)$$

$$B = \begin{bmatrix} b_{11} & m_{12} \\ m_{21} & b_{12} \end{bmatrix}, \quad ||B|| = max \begin{matrix} b_{11} + & m_{21} \\ b_{21} + & m_{12} \end{matrix}, \tag{6.3}$$

$$A = \begin{bmatrix} -m_{11} & 0 \\ 0 & -m_{22} \end{bmatrix}.$$

The logarithmic norm or the measure of any square matrix B, denoted by $\mu(B)$is defined as follows:

$$\mu(B) = \lim_{\theta \to 0+} \frac{/\!/I + \theta B/\!/ - 1}{\theta}$$

($I$ being an identity matrix of the same size as $B$).The matrix measure compatible with the norms in (6.3) is given by

$\mu(B) = \max_J [b_{jj} + \sum_{\substack{i=1 \\ i \neq J}}^{2} | b_{jJ} | ]$ $B=(b_{1J})$.

We refer to Coppel [11]or Vidyasager [31]for a number of properties of matrix measure We have shown in the appendix that the derivative of $H(Z_1,Z_2)(t)$(seeA.10)satisfies the relation

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq [-m_{11}+b_{11}+m_{21}]Z_1(t)+[-m_{22}+b_{22}+m_{12}]Z_2(t)$, and hence we have

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq \{\mu(A) + /\!/B/\!/\}\{ Z_1(t)+ Z_2(t)\}$

If we assume for the matrix norm ‖B‖ as in (6.3)that

$$\mu(A) + /\!/B/\!/ < 0 \; or \; -\mu(A) > /\!/B/\!/$$

Then from the results of the appendix (Lemmas 5 and 6) it will follow that $Z_1(t)+Z_2(t) \to 0$ as $t \to \infty$ and thus we have the following:

Lemma 6.1 Assume that the Coefficients of the linear system (6.1)satisfy the condition (6.7). Then the trivial solution of(6.1)is asymptotically stable.

When (6.7)holds an estimate for the decay rate of solutions of (6.1)can be obtained as follows A necessary and sufficient condition for the asymptotic stability of the trivial solution is that all the roots of the characteristic equation

Det $G(\lambda,T_{iJ})=0$

Have negative real parts where

$G(\lambda,T_{1J})=[\begin{matrix} \frac{\lambda+m_{11}}{0} & \frac{0}{\lambda+22} \end{matrix}] - \begin{matrix} b_{11}{}^{e-\lambda T}{}_{11} & -m_{12}{}^{e-\lambda T}{}_{12} \\ -m_{21}{}^{e-\lambda T}{}_{21} & b_{22}{}^{e-\lambda T}{}_{22} \end{matrix}$

$=\lambda I - A - B(\lambda, T_{ij})$(say),

A positive number $\mathfrak{a}$(if it exists)is said to be a decay rate of solution of (6.1) if and only if every root of (6.8)satisfies the condition

$\mathrm{Re}(\lambda)+\mathfrak{a} < 0.$

When (6.11)holds the solution of (6.1)will be of the order of $e^{-at}$ and hence the supremum of all such $\mathfrak{a} > 0$ $a$ satisfying (6.11)will be a desirable parameter to estimate However such a supremum of all such (6.11)may not be possible to determine since it has to be done directly from (6.8)-(6,10)(see also Section 3).We will obtain an estimate of $\mathfrak{a}$ (not necessarily the supremum)I (6.11)by the following procedure.

Replace $\lambda$ byz-$\mathfrak{a}$ in (6.8)and define

$G(z-\mathfrak{a}, T_{IJ}) = M(z) = (z-a)I - A - B(z-\mathfrak{a}, T_{IJ}).$

By Lemma 6.1 a Sufficient condition for all roots of M(z)=0 to have negative real parts equivalently to satisfy

$\mathrm{Re}(\lambda) < -a$

is the following :

$-\mu(aI + A) > \|B_1(a, T_{ij})\|,$

Where

$$B_1(a, T_{ij}) = \begin{pmatrix} b_{11}^{\,e-\lambda T_{11}} & -m_{12}^{\,e-\lambda T_{12}} \\ -m_{21}^{\,e-\lambda T_{21}} & b_{22}^{\,e-\lambda T_{22}} \end{pmatrix}$$

Since $\mu(\mathfrak{a} I + A) = \mathfrak{a} + \mu(A)$ we have from (6.13) that

$\mathfrak{a} + \mu(A) + \|B_I(\mathfrak{a}, T_{IJ})\| < 0$

Is a sufficient condition for Re $(\lambda) < -a$If we denote the left side of (6.15)by $\mathfrak{a} f(\mathfrak{a}, T_{Ij})$ we note that

$f(0, T_{Ij}) = \|B\| + \mu(A) < 0 \;\; if \;\; (6.7) holds;$

Al so $\partial f / \partial \mathfrak{a} > 0, f$ is continuous and monotonic in $(\mathfrak{a})$;$\to \infty as\ \mathfrak{a} \to \infty$; thus by the continuity of $f$ on $\mathfrak{a} \geq 0$, there exists an interval of the from $[0, ^{*})$such that

$f(0, T_{Ij}) < 0 \;\; for\ \mathfrak{a} \in [0, \mathfrak{a}^{*})$

And hence

$f(\mathfrak{a}, T_{Ij}) = 0$

Thus an upper bound $\mathfrak{a}^*$ of $\mathfrak{a}$ for which we have $f(\mathfrak{a}, T_{Ij}) < 0$ $\mathfrak{a} \in [0, \mathfrak{a}^*)$ can be obtained as the unique (by the monotonicity of $f$ on $\mathfrak{a}$) positive root ofmax $b_{11}^{e-\lambda T}{}_{11}$ - $m_{12}^{e-\lambda T}{}_{21}$ + $\mathfrak{a}$-min $m_{11}$ =0

$b_{22}^{e-\lambda T}{}_{22}$ - $m_{21}^{e-\lambda T}{}_{12}$ $m_{22}$

Where the maximum in (6.18) is taken over $b_{11}, b_{22}, m_{12}, m_{21}$ and $T_{ij}$ $(i, = j\ 1,2)$. A numerical procedure can be devised to solve equations of the form (6.18) in concrete cases. If one is interested in a smaller bound on the decay rate it can be obtained as the unique root of

$\|B\| e^{ar} + a + \mu(A) = 0, \tau = \max_{i,J} T_{ij}$. (6.19)

It will follow from (6.19) than an increase in T

Will reduce the value of the root of (6.19). We remark that since our estimate $a^*$ of dacay rate is based on a sufficient condition of the form (6.7), we can conclude that the solotions of (6.1) will converge at least as fast as $e^{-a*t}$ where $a^*$ is the unique root of either of (6.18) or (6.19).

## Implications

Authors have formulated a class of two systems model systems staring from the elementary principal that the rate of change of a population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) equals the creation rate minus the destruction rate. This is a departure from the conventional model building in population ecology where one usually starts from the Malthusian formulation, now known as the Kolmogorov formulation in the nonlinear context. In our analysis of the models proposed we have obtained sufficient conditions for the global existence of nonnegative solutions of the model systems (this is trivial in the Kolmogorov formulation) and the other desirable properties such as the existence of suitable steady states. The two classes of models proposed for competition and debt contain intraspecific self – regulation biologists that each systems is under the action of a self-regulating negative feedback force, opinions and arguments differ regarding the mechanism by which such an effect is realized; we refer for more details on this aspect to a collection of articles compiled by Tamarin [30].

Gopalasamy's main result is the following : if a self –regulating negative feedback mechanism (intraspecific oppugn ant) can act without time delay and if such an effect can keep an ecosystem ( a community of one or more species of populations) asymptotically stable ( at least locally about a non trivial steady state, then time delays of any magnitude in any other interaction involving positive or negative feedbacks cannot destabilize such a stable ecosystem ( in a mathematical sense); that is such delays cannot cause a hopf-type bifurcation leading to persistent and undamped fluctuations. The delay-induced bifurcations to oscillations in models have been investigated by many authors. (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Qualitatively result is not new to the literature the effects of delays on otherwise stable and unstable systems have been considered by Cushing[12]where (pages 35-39)he has shown that an otherwise unstable system can be made

less unstable by delays, Such an effect of delays on unstable systems appear to have been first proposed by Beddinton and May [2]who have also shows that delays can make stable systems less stable leading to a conclusion "time delays are not necessarily destabilizing" It is also shown by Cushing([12], pages 75-103)that while time delays cannot stabilize(in the mathematical sense)an otherwise unstable two species system in competition or mutualism time delays in there species systems can stabilize otherwise unstable systems (even in a mathematical sense)Cushing and Saleem [14]have formulated a model where delays in the positive feedback can have a stabilizing effect: Hale[18]discusses briefly the stabilizing nature of delays and given an example to show that an increase in delays can make an unstable system stable Thus for a better understanding of the effects of delays on otherwise stable systems one has to estimate the rate of decay of perturbations (or the characteristic return time )in terms of the parameters of the system especially when the delays do not cause a mathematical instability This aspect has been investigated by a procedure slightly different from that of the characteristic equation we have proposed (section6) method for the estimation of decay rate of perturbations when there are multiple delays Since many species systems with multiple delays cannot be analysed (due to complexity)by the characteristic equation our methods using comparison lead to easily verifiable sufficient conditions The sufficient conditions obtained are in the form of a diagonal dominance type; these conditions can be weakened to quasi-diagonal dominance type as in Ladde [20], Siljak[20]Anderson [1], Lewis and Anderson [21,22](Ladde's result [20] is delay dependent due to the fact that delay appears in self –regulation also).Verification of quasi –diagonal dominance conditions in arbitrary matrices (except for Metzler matrices )is an impossible task and quasi –diagonal dominance is hard to interpret.

If follows from our discussion that those confederation systems which display delay induced instability leading to permanent and undamped oscillations should have one or more of the following aspects:

(1)time delays are in intraspecific interaction terms (and possibly in other terms also): (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

(2) Intraspecific self-regulation terms are absent in density dependent from in one or more populations of the system;

(3) Intraspecific negative feedback effect are not strong enough to maintain the stability of the relevant steady state of the nondelayed system.

We can now infer how the delay induced stability leads to fluctuations; a primary mechanism of such fluctuations can be described at least theoretically as follows : in its attraction of a system to a steady state, intraspecific self-regulation fails to be effective and the system overhoots the steady state and at such higher densities intraspecific competition becomes operationally effective leading the system to understood the steady state and this process is repeated again and again resulting in undamped oscillations about the steady state. It follows from our results that if intraspecific competition is sufficiently strong, the system cannot become oscillatory even when there are time delays, so long as they do not interfere with the strength of the intraspecific self – regulation.

We are thus led to the question of existence and importance of intraspecific competition in population systems. In the introduction to an anthology of articles by several authors, Tamarin [ 30] begins with the statement "all animal populations have the reproductive capacity to increase without limit, yet none di" and in his concluding remarks writes " this leaves us with only one great concept in the area of confederation do main regulation: that all credit-debit are regulated" More specifically this is applicable to all the confederation of (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions,

financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

In his discussion on the effect of competition on the densities of populations, Nicholson ([ 25], $page$ [ 147] ) has reached the conclution "Intraspecific competition automatically regulates the severity of its action to the requirements of each denizen , provided the inherent resistance of the environment is sufficiently low to permit the species to exist." In order to look for a mechanism by which intraspecific competition plays its role of decreasing growth rate as confederation density increases, Calhoun [ 7] set up an experiment with Norway rates (Rattus norvegicus) and found that increased density leads to increased hostility which disturbs the psychology and physiology causing decreased reproductive output. Many other mechanisms by which intraspecific compition self-regulate populations have been put forward by a number of authors (see Tamarin [30]). This is true of (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

If we accept the view that intraspecific competition is really predominant in population regulation, then we can understand how and why time delays destabilize stable system leading to oscillations; time delays render intraspecific competitive regulation territorially (in case of particles we take the field in which it operates) weak so that self –regulatory forces become temporarily insufficient. If we now agree that all population models should incorporate self-regulating intraspecific competition terms in surfeit strength along with other tropic intraspecific competition terms in sufficient strength along with other atrophic interactions then it is not fair to say " time delays destabilize" and we note from our analysis that it is usually a failure in self- regulation which can cause instability rather than delays by themselves.

We will give one example to illustrate our point. Consider a credit-debit system in which the debits is self –regulating by a logistic load as follows: (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

$$\frac{dx}{dt} = r\left(x - \frac{x^2}{K}\right) - axy = b_I(x) - m_I(x,y) \quad (say),$$

$$\frac{dy}{dt} = \beta xy - \delta y = b_2(x,y) - m_2(y) - \quad (say),$$

Where $K, a, \beta, \delta$ are positive constants such that the steady state $x^*\ y^*$ exists where

$$x^* = \frac{\delta}{\beta}; \quad y^* = \frac{r}{a}\left(-\frac{\delta}{K\beta}\right) > 0.$$

In this example the predator self –regulation $(\partial m_2/\partial y)$ is not sufficiently strong in the sense

$$x^* = \left(\frac{\partial m_2}{\partial y}\right)_{(x^* y^*)} \succ \left(\frac{\partial b_2}{\partial y} + \frac{\partial m_1}{\partial y}\right)_{(x^* y^*)}$$

If in such a model there is a time delay in the credit creation rate, such a delay can lead to instability of $(x^*\ y^*)$ of setting up undamped oscillations. If however we modify the growth rate by the inclusion of a density dependent self-regulation so that credit

$$\frac{dy}{dt} = \beta\, xy - \delta y - \eta y^2$$

For some positive constant η then it is possible, depending on η, for such a modified system to remain stable without oscillations even if there is a time delay in the credit creations rate. We have another completely different justification to have $\eta y^2$ in the credit dynamics: for instance if the debit is removed from this cannot happen. While this example serves to illustrate an aspect of our result, our advocacy of the significance is not based on this example.

Gopalasamy has thus derived conclusion that "time delays need not destabilize" in models with several discrete delays. However we have provided in the appendix the necessary tools (Lemmas 3,4,5) to prove all our results when the relevant delays are all (or some) continuously distributed. Some authors (Caswell [9], May [23], Cushing [12] appear to regard continuously distributed dalays as financially/cosmologically credit wise obtained when he subjected algae to a variable nitrate environment and obtained a better fit for his data with a continuous delay. The questions of whether a discrete or continuous delay is more realistic may well depend on what process is being modeled. For instance an increased supply of some resource at some point in time may trigger reproduction at a particular moment in time later irrespective of the conditions of the resource during the intervening period. If the duration of increased resource supply necessary to elicit such a response is short compared to the intervening period before the reproduction occurs, then a discrete time delay may be appropriate. In a large number of cases of animal. Note fungus algae are prey. Here it shall be an account which loses the energy, amount, or matter for that matter. Same result holds good for (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Populations of particles, atoms, molecules, propositions, systems types chosen the past dietary and nutritional history of an animal over a long period plays an influential role in determining the current bahaviour of the dietary and nutritional history of an animal over a long period plays an influential role in determining the current behavior of the animal and in such cases continuous time delays will be appropriate.

In conclusion the following comments are appropriate. The time delays since the current creation rate is usually due to a contribution of those mature animals living a few units of time earlier depending on the type of animal or organism; the time delays appearing in interspecific interactions can be due to the fact that in order to successfully interact and compete interferentially a developmental maturity may be necessary. In the case of predation, delays in the birth rate of prey can be due to delayed maturation while that in the birth rate of a predator is due to gestation delay with delayed fecundity; the delays in the predation rate are again are again as before due to the maturity needed for a predator to successfully organize and succeed in an attack.

In our analysis we have assumed throughout a temporally uniform environment neglecting dispersal of populations and their age structure. We expect our results to hold under suitable conditions both in temporally periodic environments with the steady state of a constant environment being replaced by a periodic solution

(nonstationary steady state) and when the effects of dispersion and age structures are considered. It is worthwhile to modify our purely ecological, cosmological, financial, propositional, problematic systems models so as to provide them with a genetic foundation and examine the effects of time delays since forces of selection act with some time delay. Such a generalization will definitely lead to a better understanding of "how confederations regulate themselves".

## Appendix

First we derive a comparison result.

Lemma 1. Consider the delay differential system

$$\left.\begin{aligned}\frac{dY_1(t)}{dt} &= -m_{11}\ Y_1(t) + b_{11}Y_1(t-\tau_{11}) - m_{12}\ Y_2(t-\tau_{12}),\\ \frac{dY_2(t)}{dt} &= -m_{22}Y_2(t) - m_{21}Y_1(t-\tau_{21}) + b_{22}\ Y_2(t-\tau_{22}),\end{aligned}\right\} (A.1),$$

Where the $m_{ij}$, $b_{ij}$ $(i.j = 1,2)$ are positive while $\tau_{ij}$ $(i.j = 1,2)$ nonnegative with $\tau = \max \tau_{ij}$, $i,j = 1,2$. Then

$$\left.\begin{aligned}\frac{d|Y_1(t)|}{dt} &\le -m_{11}\,|Y_1(t)| + b_{11}|Y_1(t-\tau_{11})| + m_{12}\ Y_2(t-\tau_{12})|,\\ \frac{d|Y_1(t)|}{dt} &\le -m_{22}\,|Y_2(t)| + m_{21}|Y_1(t-\tau_{21})| + b_{22}\ |Y_2(t-\tau_{21})|,\end{aligned}\right\} (A.2)$$

Where

$$\frac{d}{dt}\,|\,Y_1(t)\,| = \left\{\begin{aligned}&\frac{dY_1(t)}{dt}\ if\ \ Y_1(t) \ge 0,\\ &\frac{dY_1(t)}{dt}\ if\ \ Y_1(t) < 0,\end{aligned}\right\}\ i = 1,2. \qquad (A.3)$$

The proof is by direct verification.

LEMMA 2. Consider now (A.2) and the following system

$$\left.\begin{aligned}\frac{dZ_1(t)}{dt} &= -m_{11}\ Z_1(t) + b_{11}Z_1(t-\tau_{11}) + m_{12}\ Z_2(t-\tau_{12}),\\ \frac{dZ_2(t)}{dt} &= -m_{22}Z_2(t) + m_{21}Z_1(t-\tau_{21}) + b_{22}\ Z_2(t-\tau_{22}),\end{aligned}\right\} \qquad (A.4)$$

Together with the initial conditions

$$Z_1(s) > |Y_1(s)|,\ \ s \in [-\tau, 0],\ i = 1,2. \qquad (A.5)$$

Then

$$|Y_1(t)| < Z_1(t) \text{ for } t \ge 0 \quad \text{and} \quad i = 1,2. \qquad (A.6)$$

PROOF. Let $U_1(t) = Z_i(t) - |Y_1(s)|$, $i = 1,2.$ $t \geq -\tau.$ Then we have from (A.2) and (A.4),

$$\left.\begin{aligned} \frac{dU_1(t)}{dt} &\geq -m_{11}\ U_1(t) + b_{11}U_1(t-\tau_{11}) + m_{12}\ U_2(t-\tau_{12}), \\ \frac{dU_2(t)}{dt} &\geq -m_{22}U_2(t) + m_{21}U_1(t-\tau_{21}) + b_{22}\ U_2(t-\tau_{22}), \end{aligned}\right\} t > 0, \quad (A.7)$$

With

$$U_1(s) = Z_i(s) - |Y_i(s)| > 0, \quad i = 1,2, s \in [-\tau, 0]. \quad (A.8)$$

Consider (A.7)-(A.8) first on the interval $[\tau, 0]$ : It will follow from (A.7) that $U_1(t) \geq U_1(0) \exp [-m_{11}t]$

$$+\int_0^t [b_{11}U_1(s-\tau_{11}) + m_{12}\ U_2(s-\tau_{12})] \exp [-m_{11}(t-s)ds > 0,] \quad (A.9)$$

For all $t \in [0, \tau^*]$ since $U_i(s) > 0$ on $[-\tau^* 0]$ where $\tau^*$ = positive minimum of $(\tau_{11,}\tau_{12,}\tau_{21,}\tau_{22,})$. Similarly $U_2(t) > 0$ for $t \in [0, \tau^*]$. Now by a similar argument it will follow that $U_1(t) > 0, U_2(t) > 0$ for $t \in [\tau^*, 2\tau^*]$. Repeating the above procedure we obtain the result (A.6).

It follows from Lemmas 1 and 2 that in order to prove theorem 3.1 it is enough to prove that $Z_i(t) \to 0$ as $t \to \infty$ for $i = 1,2$ . which we will do as follows.

PROOF OF THEOREM 3.1 Define a Lyapunov-like function

$$H(Z_1,Z_2)(t)=Z_1(t)+Z_2(t)$$

$$+\int_{1-T}^t b_{11}\, Z_1(s)ds + \int_{1-T}^t m_{12}\, Z_2(s)ds$$

$$+\int_{1-T}^t m_{21}\, Z_1(s)ds + \int_{1-T}^t b_{22}\, Z_{22}(s)ds$$

Since each term of H in (A.10)is nonnegative for t>0, if we show that $H \to 0$ as $t \to \infty$ it will follow from Lemma 2 that

$$0 \leq |\, Y_{1,}(t)\,| \leq Z_1(t) \to 0 \text{ as } t \to \infty,$$

Implying that the perturbations $Y_{1,}$, $Y_2$, governed by (A.1)approach zero as $t \to \infty$, and this will prove our result Calculate the rate of change of H along the solutions of (A.4)so that

$$\frac{dH}{dt} = \frac{\partial H}{\partial t} + \frac{\partial H}{\partial Z_1}\frac{dZ_1}{dt} + \frac{\partial H}{\partial Z_2}\frac{dZ_2}{dt}$$

$$=-[Z_1(t)\{m_{11} - (b_{11} + m_{12})\} + Z_2(t)\{m_{22} - (b_{22}m_{12})\}] < 0$$

On using (2.10)since $Z_1(t),Z_2(t)$ are positive for all t>0. Thus $H(Z_1, Z_2)(t) \to 0$ as $t \to \infty$ since H is bounded below by zero: it follows from Lemma 6 below that $Z_1(t) \to 0, Z_2(t) \to 0$ as $t \to \infty$ the proof is complete.

The following set of results are intended to be a guide to prove result similar to the above in the case of continuously distributed delays An analogue of Lemma I is the following.

LEMMA 3 Let $K_{ij}(s)$,s∈[0,∞), $i,j = 1,2$,denote nonnegative continuous functions integrable on [0,∞) normalised such that

$\int_0^\infty sk_{ij}(s)ds< \infty$; $\int_0^\infty k_{ij}(s)ds=1$, I,j=1,2.

Consider the intergrodifferential system

$$\frac{dY_1(t)}{dt} = -m_{11}Y_1(t) + b_{11}\int_0^\infty k_{ij}(s)\, Y_1(t-s)ds$$

$$-m_{12}\int_0^\infty k_{12}(s)\, Y_2(t-s)ds$$

$$\frac{dY_2(t)}{dt} = -m_{22}Y_2(t) + m_{21}\int_0^\infty k_{21}(s)\, Y_1(t-s)ds$$

$$+b_{22}\int_0^\infty k_{22}(s)\, Y_2(t-s)ds$$

*Where* $m_{ij}, b_{ij}$ $(i,j = 1,2)$ *are positive constants Then*

$$\frac{d|Y_1(t)|}{dt} \leq -m_{11}|\,Y_1(t)\,| + b_{11}\int_0^\infty k_{11}(s)|Y_1(t-s)|ds,$$

$$+m_{12}\int_0^\infty k_{12}(s)|\,Y_2(t-s)|ds,$$

$$\frac{d|Y_2(t)|}{dt} \leq -m_{22}|\,Y_2(t)\,| + m_{21}\int_0^\infty k_{21}(s)|Y_1(t-s)|ds,$$

$$+b_{22}\int_0^\infty k_{22}(s)|\,Y_2(t-s)|ds,$$

*Where* $d|Y_1(t)||/dt$ *is defined by* (A.3).

Proof is again by direct verification.

The proof of the following lemma is similar to that of Lemma 2 above.

Lemma 4. *Consider (A.14) and the following intergrodifferential system:*

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1(t) + b_{11}\int_0^\infty k_{11}(s)Z_1(t-s)ds$$

$$+m_{12}\int_0^\infty k_{12}(s)Z_2(t-s)ds,$$

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1(t) + b_{11}\int_0^\infty k_{11}(s)Z_1(t-s)ds$$

$$+b_{22}\int_0^\infty k_{22}(s)Z_2(t-s)ds,$$

Together with the initial conditions

$$\infty > Z_i(s) > |Y_i(s)| for\ s \in (-\infty, 0).$$

(A.16)

Then $|Y_i(t)| < Z_i(t) for\ t \geq 0, i = 1,2.$

The following lemma will be useful to prove a result analogous to that of Theorem 3.1 corresponding to a class of continuously distributed time delays.

LEMMA 5. Let $Z_1(t), Z_2(t)$ be the solutions of (A.15)- (A.16) where $(Y_1\ Y_2)$ corresponds to an arbitrary but fixed solutions of (A.13). Then as a consequence of Lemma & it will follow that

$$0 \leq |Y_i(t)| \ Z_i(t) \to \infty \leq as\ t \to \infty, \qquad i = 1,2,$$

(A.17)

Whenever

$$\begin{cases} M_{11} > b_{11} + m_{21}, \\ M_{22} > b_{22} + m_{12} \end{cases}$$

PROOF. Define a Lyapunov –like function H where

$$H(Z_1,Z_2)(t) = Z_1(t) + Z_2(t) + b_{11}\int_0^\infty k_{11}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$$

$$+ m_{12}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$$

$$+ m_{21}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$$

$$+ b_{22}\int_0^\infty k_{22}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$$

Where $(Z_1,Z_2)$is any solution of (A.15)-(A.16) for arbitrary $(Y_1,Y_2)$defined on $(-\infty, 0]$ Computing the rate of change of H in (A.19)along the solution of (A.15)we derive that (on using (A.12))

$$\left(\frac{dH(Z_1Z_2(t)}{dt} = [Z_1(t)\{m_{11} - (b_{11} + m_{21})\} + Z_2(t)\{m_{22} - (b_{22} + m_{12})\}]\right.$$

$< 0$(using (A.18)) for $Z_1(t) > 0\ and\ Z_2(t) > 0.$

By Lemma 6 below the result will follow

Lemma 6. Let H(t)=H$(Z_1, Z_2)(t)$ be as defined in (A.10) Suppose

$$m_{11} > b_{11} + m_{21} \text{ and } m_{22} > b_{22} |m_{12}$$

Then

$$Z_1(t) + Z_2(t) \to 0 \text{ast} \to \infty.$$

PROOF. Suppose (A.21) dose not hold; then there exists at least one sequence

$\{t_n\} \to \infty$ as n $\to \infty$ such that

$$[m_{11} - (b_{11} + m_{21})]\, Z_1(t_n) + [m_{22} - (b_{22} + m_{12})]Z_2(t_n) > \varepsilon \qquad (A.22)$$

For some sufficiently small positive number $\varepsilon$ and hence

$$\frac{dH(t_n)}{dt} < -\varepsilon, n = 1,2,3, \ldots.$$

By the continuity and differentiability of H with respect to t,

$$\mathrm{H}(t_n - \varepsilon) - \mathrm{H}(t_n) = \frac{dH(t^*)}{dt}(-\varepsilon), \quad t_n - \varepsilon < t^* < t_n$$

If $\varepsilon$ is small we will have

$$= \frac{dH(t^*)}{dt}\varepsilon/2$$

So that

$$\mathrm{H}(t_n) \leq \mathrm{H}(t_n - \varepsilon) - \varepsilon^2/2$$

$$\leq H(t_{n-1}) - \varepsilon^2/2$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2)$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2):$$

$$\leq H(t_0) - \mathrm{n}(\varepsilon^2/2): \mathrm{t}_0 = 0.$$

(A.12) and Lemma 5 will imply that $H(t_0) < \infty$ for any $t_0 > 0$. If n is sufficiently large it will follow that $H(t_n) < 0$ since $H(t_0) < \infty$, which contradicts the nonnegativity of H for all $t \geq 0$. Hence $Z_1(t) + Z_2(t) \to 0$ as $t \to \infty$ implying the asymptotic stability of the trivial solution of (A.1)(since $Z_1$ dominates $y_1$)

We note the result of Lemma 5 and 6 will remain valid if (A.18) is replaced by the condition $\mu(A) + \|B\| < 0$

Where A and B are defined as in (6.3).

We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities,

apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems

**Gopalsamy**

# Model For Innocuously Benign Delay In A Periodic System

**Abstract**

Sufficient conditions are obtained for the existence of a unique linearly denomination stable positive periodic solution of a system model of two competitions in a periodic environment with time lags in interspecific interactions. It is shown that if the self- regulating intraspecific interaction effects are strong enough and act without time delays then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a periodic domination. An example in finance will make things clear. Corporate finance for one big plant generates lot of accounts of the workers whose salary shall be disbursed to the debit of cash credit account. Here Cash credit account is a resource as also the one to which debits are raised, a la predator. Some of the accounts from which the credits come are demand bills purchase, usance bills discount, vostro and nostro accounts in respect of foreign bills and Letters of credit. In economics inventories are used to produce output, output is sold to produce income which might be saved in part and recycled in investment. Cosmologically take an example of a Black hole which gormandizes matter and radiates energy. Notwithstanding information paradox, one thing is clear that matter is being created in to radiation. We show how the model could be applied to axiomatic predications, postulation alcovishness, phenomenological methodologies, propositional statement, and other problems in cosmology, philosophy, and even in Green House effects. We study the sine qua non properties of such a system. In fact these models can be applied to such unsolved problems like Yang mills Theory and mass gap, and other problems. As we apply the model, we shall go on dilating the position status and the applicability of the model. We shall not touch the electrical engineering problems which might rise questions like what is charge of a Bank. We may say when a customer enters a Bank; we can say the Bank is charged. These factors are still relativistic and we shall make an in-depth study as we progress. Gopalsamy's research was published in Australian Journal of Mathematical Society vide: B 25349=365 It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 16. Introduction

At the outset we assume that the system satisfies Haimovici is caution. The purpose of this article is to establish a set of sufficient conditions for the existence of a unique linearly asymptotically stable positive periodic solution for a nonautonomous system of the form

$$\frac{dx(t)}{dt} = x(t)f_1\big(t, x(t), y(t-\tau_1)\big),$$

**(1.1)**

$$\frac{dy(t)}{dt} = y(t)f_2\,\big(t, x(t-\tau_2), y(t)\big).$$

In the context of confederation dynamics a model system of the form (1.1) means that the interspecific interactions occur with time delays $\tau_1$ 0, $\tau_2$ 0 in a periodic environment, if $f_1$ *and* $f_2$ are assumed to be periodic in $t$ with a common period $\omega$ 0. Note the interactions are inter specific.

Although model systems in periodic domains have been investigated by a large number of authors (for some references see Gopalsamy (4], [5]), periodically varying environments in time delayed model systems of two or more interacting species have not been investigated. The need for Multitype system models in periodic environments and time delayed interactions are evident from the extensive literature of that type relating to single species. While it is known that periodic ecosystems can support more species than uniform environments (Gopalsamy [4]), time delays might destabilize the system leading to a possible loss of one or more species. Furthermore, periodic environments in population dynamics have some relevance to natural selection and evolution; we will consider this aspect briefly at the end.

The principal finding of the following analysis of (1.1) is that, under suitable hypotheses on $f_1$ *and* $f_2$, the system (1.1) has a linearly asymptotically stable unique positive periodic solution for all nonnegative delays $\tau_1, \tau_2$ (while our numerical experiments suggest that the periodic solution of (1.1) is globally asymptotically stable and is independent of the delay parameters $\tau_1, \tau_2$ it is not known how one can prove such a strong result). Since time delays are usually associated with possible destabilization of otherwise stable systems, we can describe the time delays appearing on (1.1) as “harmless”.

It has been recently shown by the author (Gopalsamy [6]) that if the self –regulating negative feedback in each denomination is strong enough and acts without time delays, then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a temporally uniform domain. The result of this paper is thus an extension to a periodic environment of the "harmless" nature of interspecific interaction delays in a temporally constant environment mentioned above.

## 17. Existence and uniqueness of periodic solutions

We will now formulate the problem of existence of periodic solutions of the periodic delay differential system

$$\frac{dx_1(t)}{dt} = x_1(t) f_1\big(t, x_1(t), x_2(t-\tau_{12})\big),$$

$$t > 0,$$

**(2.1)**

$$\frac{dx_2(t)}{dt} = x_2(t) f_2\big(t, x_1(t-\tau_{21}), x_2(t)\big),$$

(Where $\tau_{12}, \tau_{21}$ are nonnegative constants) in a form suitable for the application of fixed point theory of positive operators. We make the following assumptions regarding $f_1$ and $f_2$ in (2.1).

($H_1$) The partial (Frechet) derivatives of $f_1, f_2$ satisfy the following:

$$\frac{\partial f_i}{\partial x_j} < 0, \qquad i, j = 1, 2 \text{ For all } t > -\tau \; where \; \tau = \max(\tau_{12}, \tau_{21}). \qquad \textbf{(2.2)}$$

($H_2$) $f_1(t, x_1, x_2) \equiv f_1(t + \omega, x_1, x_2)$ for all $t > -\tau, x_1 \; \square 0, x_2 \; \square 0$. **(2.3)**

($H_3$) There exist positive constants $b_{il}$, $b_i^u$, $a_{il}$, $a_i^u$ ($l = 1, 2$) such that

$$b_{il} - a_{il} x_i < f_i(t, x_1, x_2) < b_i^u - a_i^u x_i,$$

$$b_{il}/a_{il} < b_i^u / a_i^u, \qquad \textbf{(2.4)}$$

For $i$= 1, 2 and $(x_1, x_2) \in R_+ \times R_+$. If $\tau = 0$ then one can show that (2.1)-(2.4) will imply that the nondelayed periodic system (2.1) will have an asymptotically stable strictly positive periodic solution (for details see Gopalsamy [5]); this result will become a special case of what we obtain in the following.

We will use the following notation:

$$B_\omega^\tau = \{h: [-\tau, \infty) \to R, h \; is \; continuous, h(t + \omega) = h(t), t \; \square -\tau\},$$

$$B = B_\omega^\tau \times B_\omega^\tau,$$

$$K_1 = \{ h \in B_\omega^\tau | h(t) \geq 0, t \in [-\tau, \infty)\},$$

$$K = K_1 \times K_1,$$

$$||h|| = \left\{ \sup_{t \in [0,\omega],} | h(t)|, h \in B_\omega^\tau \right\}$$

$$||x_1, x_2|| = \max\{||x_1||, ||x_2||, (x_1, x_2) \in B\},$$

$$G_1 = \{(x_1, x_2) \in K, ||(x_1, x_2)|| \leq b_l = \min(b_{1l} / a_{1l}, b_{2l}/a_{2l})\}$$

$$G_2 = \left\{(x_1, x_2) \in K, ||(x_1, x_2)|| \geq b^u = \max\left(\frac{b_1^u}{a_1^u}, \frac{b_2^u}{a_2^u}\right)\right\},$$

$\bar{G} =$ closure of $G_2$ *in* $B$ and

$\frac{\bar{G}_2}{G_1} =$ complement of $G_1 in\ \bar{G}_2$; $G_1 \subset G_2$ by **(2.4)**

The cone $K = K_1 \times K_1$ induces a partial order in $B$ and we will write $(x_1, x_2) \geq (y_1, y_2)$ if and only if $x_1 - y_1 \in K_1$. we define an operator $N$ on $B$ as follows

$$N(x_1, x_2)(t) = \begin{matrix} N_1(t, x_1, x_2) \\ N_2(t, x_1, x_2) \end{matrix} \quad \text{for t> } 0, \qquad \textbf{(2.5)}$$

Where

$$N_1(t, x_1, x_2) = \int_0^\omega G_{11}(m, \omega, t-s) g_1\left(m, s, x_1(x) x_2(s - \tau_{12})\right) ds$$

$$+ \int_0^1 G_{12}(m, 0, t-s) g_1\left(m, s, x_1(s), x_2(s - \tau_{12})\right) ds,$$

$$N_2(t, x_1, x_2) = \int_0^\omega G_{22}(m, \omega, t-s) g_2\left(m, s, x_1(s - \tau_{21}), x_2(s)\right) ds \qquad \textbf{(2.6)}$$

$$+ \int_0^1 G_{21}(m, 0, t-s) g_2\left(m, s, x_1(s - \tau_{21}), x_2(s)\right) ds,$$

For $t > 0$, with

$$g_1\left(m, s, x_1(s), x_2(s - \tau_{12})\right) = m x_1(s) + x_1(s) f_1\left(s, x_1(s), x_2(s - \tau_{12})\right),$$

**(2.7)**

$$g_2\big(m,s,x_1(s-\tau_{21}),x_2(s)\big) = mx_2(s) + x_2(s) f_2\big(s,x_1(s-\tau_{21})x_2,(s)\big),$$

$$G_{ll}(m,\omega,t-s) = e^{-m\omega}(1-e^{-m\omega})^{-1}e^{-m(t-s)}, \qquad i=1,2,$$

**(2.8)**

$$G_{ij}(m,0,t-s) = e^{-m(t-s)}, i \neq j, i,j = 1,2,$$

$$N(x_1,x_2)(t) = \begin{cases} N_1(n\omega+t,x_1,x_2) \\ N_2(n\omega+t,x_1,x_2) \end{cases} for\ t \in [-\tau,0], \qquad \textbf{(2.9)}$$

*m* being a suitably selected positive constant (more will be said below) while *n* is a positive constant such that $n\omega > \tau$.

We note immediately the following from the definition of $N$: *if* $(x_1,x_2) \in \bar{G}/G_1$ then $N(x_1,x_2)(t)$ and if $N(x_1,x_2)$ is periodic $\omega$ for $t$ then $N(x_1,x_2)$ is periodic in *t* with period $\omega$ for $t \geq -\tau$. The existence of at least one positive periodic solution of (2.1)-(2.4) is established by means of a few lemmas and a result of Schmitt [9].

Lemma 1. *A necessary and sufficient condition for* $(x_1(t),x_2(t))$, $t \geq -\tau$ *to be a periodic solution of (2.1)-(2.4) of period* $\omega$ *is that* $(x_1,x_2)$ *be a fixed point of N such that*

$$N(x_1,x_2)(t) = \big(x_1(t),x_2(t)\big) \quad and\ (x_1,x_2) \in K \cap \left(\frac{\bar{G}}{G_1}\right).$$

PROOF. The details of proof are similar to those of Lemma 2 of Busenberg and Cooke [1]. We first verify the continuity of $N(x_1,x_2)$ assuming that $N(x_1,x_2)(t)$ is periodic in *t* of period $\omega$. The continuity of $N(x_1,x_2)$ for $t > 0$ follows from the definition in (2.5)-(2.8) WHICH ALSO PROVIDES $\lim_{t\to 0+} N(x_1,x_2)(t)$. It will follow from (2.9) and the assumed periodicity of $N(x_1,x_2)$ *that* $N(x_1,x_2)$ is continuous

For $t \geq -\tau$. Now we will verify the periodicity of $N(x_1,x_2)$ in $t$. A simple computation leads to

$$N_1(t+w,x_1,x_2) - N_1(t,x_1,x_2) = e^{-mt}F_1(t),$$

**(2.10)**

Where

$$F_1(t) = \int_0^{\omega} G_{11}\,(m, \omega, t-s\,)g_1\big(m, s, x_1(s), x_2(s-\tau_{12})\big)ds$$

$$+\int_0^{\omega} G_{12}(m, 0, t-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$+\int_{\omega}^{\omega+t} G_{12}(m, 0, \omega-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$-\int_0^{\omega} G_{11}(m, \omega-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$+\int_0^{t} G_{12}(m, 0, -s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

**(2.11)**

Which on simplification using the periodicity of $f_1$ shows that $F_1(t) \equiv 0$ whenever $x_1, x_2, \in B$, thus implying the periodicity of $N_1$ in $t$ of period$\omega$. A similar argument for $N_2$ proves that $(x_1, x_2) \in B \Rightarrow N\,(x_1, x_2) \in B$.

Now suppose $(x_1, x_2)$ is a fixed point of *N and let* $(x_1, x_2)\ \in K \cap (\bar{G}/G_1\,)$. Then it is easily seen that $(x_1, x_2)$ is a solution of (2.1) and since $(x_1, x_2)\ \in K \cap (\bar{G}/G_1\,)$ it will follow that $(x_1, x_2)$ is a periodic solution of (2.1).

Conversely, let $(x_1, x_2)$ be a periodic solution of period $\omega$ of (2.1). Then $(x_1, x_2)$ is necessarily a periodic solution of

$$\frac{dx_1(t)}{dt} + mx_1(t) = g_1\,\big(m, t, x_1\,(t), x_2(t-\tau_{12})\big)$$

**(2.12)**

$$\frac{dx_2(t)}{dt} + mx_2(t) = g_2\,\big(m, t, x_1\,(t-\tau_{21}), x_2(t)\big),$$

For any arbitrary real constant *m* and hence for any fixed constant $m > 0$. It follows from (2.12) that any periodic solution of (2.1) and hence of (2.12) is a fixed point of $N(x_1, x_2)$. We claim that such a fixed point is in $K \cap (\bar{G}_2/G_1)$. Suppose $(x_1, x_2) \notin K \cap (\bar{G}_2/G_1)$. the following possibilities can arise:

(xvi) Either $x_1, x_2$ or both $x_1$ *and* $x_2$ may become negative for some value of *t* and it will follow from (2.1) that subsequent value of $x_1, x_2$ or $x_1$ *and* $x_2$ will be negative and hence from (2.4) and (2.1) their derivatives will be negative for all subsequent *t* , implying that $x_1, x_2$ or $x_1$ *and* $x_2$ cannot be periodic.

(xvii) While $(x_1, x_2) \in K \subset B$ there may be a value of *t*, say $t^*$, such that $x_1(t^*) > b_1^u/a_1^u$ or $x_2(t^*) > b_1^u/a_1^u$ and in such a case $x_1(t) or\ x_2(t)$ cannot be periodic since respectively $dx_1(t^*)/dt < 0\ or\ dx_2(t^*)/dt < 0,$ which will imply that the values of $x_1(t^*)\ or\ x_2(t^*)$ cannot be repeated from below.

(xviii) Similarly if there is a $t^*$ such that $0 < x_1(t^*) < b_{1l}/a_{1l}\ or\ 0 < x_2(t^*) < b_{2l}/a_{2l}$ then such solutions cannot be periodic since $dx_1(t^*)dt > 0$ or $dx_2(t^*)dt > 0$, which will mean that the values of $x_1(t^*)$ or $x_2(t^*)$ cannot be repeated from above.

Thus every periodic solution of (2.1)-(2.4) is a fixed point of *N* in $K \cap \left(\frac{G_2}{G_1}\right) \subset B$.

**LEMMA 2.** *Let N be as in (2.5)-(2.6). Suppose* $\lambda \in (0,1) and\ \lambda N(x_1, x_2) = (x_1, x_2)$ *for some* $((x_1, x_2) \in K - \{0\}$. *then there exist* $\alpha_1^* > 0, \alpha_2^* > 0$ *such that*

$$0 < x_1(t) \leq \alpha_1^* < b_1^u / a_1^u,$$

$$0 < x_2(t) \leq \alpha_2^* < b_2^u / a_2^u,$$

*for* $t \in [-\tau, \infty)$.

PROOF. If $\lambda N(x_1, x_2) = (x_1, x_2)$ then from the definition of N we get for all $t > 0$ and $(x_1, x_2) \in K - \{0\}$,

$$\frac{dx_1(t)}{dt} = \lambda\, x_1\,(t) f_2(t, x_1(t), x_2\,(t-\tau_{12})) + x_1\,(t) m(\lambda - 1)$$

**(2.13)**

$$\frac{dx_2(t)}{dt} = \lambda\, x_2\,(t) f_2(t, x_1(t-\tau_{21}), x_2\,(t)) + x_2\,(t) m(\lambda - 1)$$

For all $t \geq 0$ such that $x_1(t) \geq b_1^u/a_1^u$ *or* $x_2(t) \geq b_2^u/a_2^u$ we have $\frac{dx_1}{dt} < 0$ *or* $\frac{dx_2}{dt} < 0$; but these are impossible since $x_1, x_2$ are periodic in t. Thus by the continuity and periodicity of $x_1, x_2$ we have $x_1(t) \leq x_1^* < b_1^u/a_1^u$, $x_2(t) \leq x_2^* < b_2^u/a_2^u$ for all $t \geq 0$; by the definition and periodicity the same will follow for all $t \geq -\tau$. A similar argument holds when $x_1(t) < 0$ or $x_2(t) < 0$, noting that in such a case $dx_1/dt > 0$ or $dx_2/dt > 0$, again contradicting the periodicity of $(x_1, x_2)$ in *t*.

LEMMA 3. *The operator N maps* $K \cap (G_2/G_1)$ *in to K and is completely continuous*

PROOF. From the definition of N and (2.4) we derive that

$$N_1(t, x_1, x_2) \geq \int_0^\omega G_{11}\,(m, \omega, t-s) x_1(s)[m + b_{1l} - a_{1l} x_1(s)] ds$$

$$+ \int_0^\omega G_{12}\,(m, 0, t-s) x_1(s)[m + b_{1l} - a_{1l} x_1(s)] ds.$$

**(2.14)**

For $(x_1, x_2) \in K \cap (G_2/G_1)$, we have $0 < x_1 < b_1^u/a_1^u$ and hence we get $N_1(t, x_1, x_2) > 0$ for all $t \geq 0$ by a proper choice of the positive constant say $m_1$.

We make a similar choice say $m_2$ for *m* for the case $N_1(t, x_1, x_2) \geq 0$ and then choose *m* = max $(m_1, m_2)$. Thus we have

$$N: K \cap (G_2/G_1) \to K \qquad for\ t \geq 0,$$

And by the periodicity of N we have the same for $t \in -\tau$.

To prove the complete continuity of N we proceed as follows. We know that if $(x_1, x_2) \in B$ then $N(x_1, x_2) \in B$ and so we can restrict $t$ to the set $[0, n\,\omega], n\omega > r$. We write $N = L \circ Q$ where $L$ is the linear operator, $L : B \rightarrow B$, $L(u_1, u_2)(t)$

$$= \int_0^{\omega} G_{11}(m, \omega, t-s) \begin{bmatrix} u_1 & (s) \\ u_2 & (s) \end{bmatrix} ds, + \int_0^{t} G_{12}(m, 0, t-s) \begin{bmatrix} u_1 & (s) \\ u_2 & (s) \end{bmatrix} ds$$

**(2.15)**

And Q is the nonlinear operator

$$Q(u_1, u_2)(t) = \begin{matrix} g_1(m, t, u_1(t), u(t-\tau_{12})) \\ g_2(m, t, u_1(t-\tau_{21}), u_2(t)) \end{matrix} for\ \ t > 0,$$

**(2.16)**

$$Q(u_1, u_2)(t) = Q\,(u_1, u_2)(n\omega + t) \quad for\ t \in [-\tau, 0].$$

It can be found that

$$||L(u_1, u_2)(t)|| \leq max \left\{ \frac{m+b_1^u}{m}, \frac{m+b_2^u}{m} \right\} ||u_1, u_2||,$$

**(2.17)**

$$\left\| \frac{d}{dt} L(u_1, u_2) \right\| \leq max \left\{ \frac{m+b_1^u}{m}, \frac{m+b_2^u}{m} \right\} ||u_1, u_2||.$$

The complete continuity of the linear map $L$ follows from (2.17). Now since Q is continuous and maps bounded sets into bounded sets it will follow that $L \circ Q$ is completely continuous.

To prove the existence of periodic solution of (2.1)-(2.4) we will use the following result of Schmitt [9].

**THEOREM (Schmitt).** *Let $H_1$ and $H_2$ be nonempty bounded open neighborhoods of the origin in a Banach space B with a positive cone K such that $H_1 \subset H_2$ and let N: $K \cap (H_2 / H_1) \to K$ be completely continuous and which satisfies the following conditions:*

***(k)*** *There exists $(z_1, z_2) \in K \subset B, \|(z_1, z_2)\| = 1$ and a $\mu$ satisfying*

$$\mu > sup\{\|(x_1, x_2) - N(y_1, y_2)\|, (x_1, x_2) \in K \cap H_1, (y_1, y_2) \in K \cap \partial H_2\}$$ *such that all solution* $(x_1, x_2) \in K \cap (H_2/H_1)$ *of* $(x_1, x_2) = N(x_1, x_2) + \lambda(z_1, z_2) with\ 0 < \lambda < \mu\ satisfy\ (x_1, x_2) \notin \partial H_1.$

*(l) All solution* $(x_1, x_2) \in K \cap (H_2/H_1)\ of\ (x_1, x_2) = \lambda N(x_1, x_2), 0 < \lambda < 1\ satisfy\ (x_1, x_2) \notin \partial H_2.$

*Then N has at least one fixed point in* $K \cap (H_2/H_1)$.

We can now present our existence result.

**THEOREM 1.** *Assume that* $f_1, f_2$, *of* (2.1) *satisfy* (2.2)-(2.4). *then there exists at least one positive periodic solution of* (2.1)-(2.4).

PROOF. By Lemma 1, it is enough to show that the operator $N: K \cap (G_2/G_1) \to K \subset B$ has at least one fixed point in $G_2/G_1$ and we will do this by applying the result of Schmitt [9]. First we let

$$H_2 = \{(x_1, x_2) \in B | \ ||x_1, x_2|| < \max(b_1^u/a_1^u, b_2^u/a_2^u)\} \equiv G_2,$$

$$H_1 = \{(x_1, x_2) \in B | \ ||x_1, x_2|| < \min(b_{1l}/a_{1l}, b_{1l}/a_{1l})\} \equiv G_1.$$

By Lemma 3, $N: K \cap (H_2/H_1) \to K$ is completely continuous. Choose $\mu = \infty$ and $(z_1, z_2) \equiv (1.1)$. Then all solution $(x_1, x_2) \in K \cap (H_2/H_1)$ of $(x_1, x_2) = N(x_1, x_2) + \lambda(z_1, z_2), 0 < \lambda < \infty$, satisfy (2.1) and hence by their periodicity $(x_1, x_2) \notin \partial H_1$. The condition (a) of Schmitt's theorem is satisfied since $\mu = \infty$ by choice and $(x_1, x_2)$ and $(z_1, z_2)$ are periodic. By Lemma 2 the condition (b) of Schmitt's theorem is satisfied. Thus there exists at least

one fixed point of N in $K \cap (H_2/H_1)$ and hence there exists at least one positive periodic solution of (2.1)-(2.4) and the proof is complete.

We remark that so far no restriction on the magnitudes of the delays $\tau_{12}, \tau_{21}$ has been imposed except that they be nonnegative.

## 18. Uniqueness of the periodic solution

Applications of fixed point theorems to prove the existence of solutions of nonlinear problems usually suffer from the drawback that the uniqueness of solutions has to be proved often by other methods. We will use the following result of Smith [11] which is a slight variation of his Theorem 3.5 (see also Busenberg and Cooke [1]).

THEOREM (Smith). *Let N be an operator on a cone K in the Banach space B and let* $K^* = \{(x_1, x_2) \in K \cap B, 0 \leq$ *x1,x2≤C∗ where C∗some positive constant is. Assume the following:*

(p) If $(x_1, x_2) \in K^* - \{(0,0)\}$ and $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

(q) N is monotonic nondecreasing on $K^*$.

(r) If $(x_1, x_2)$ is in the interior of $K^*$ and if $0 < \lambda < 1$ then there exists an $\eta > 0$ depending on $(x_1, x_2)$ and $\lambda$ such that $N(\lambda x_1, \lambda x_2) \geq \lambda(1+\eta) N(x_1, x_2)$. Then N has at most one fixed point different from $(0,0)$ in $K^* \subset B$.

The cone K= $K_1 \cap K_1$ where $K_1 = \{h \in B_\omega^\tau : h(t) \geq 0, t \in [-\tau, \infty)\}$ is slightly too wide in that solutions of the form $\{x_1(t), 0\}$ and $\{0, x_2(t)\}$ for (2.1) are also possible. We choose a narrower cone $K^*$ in $B$ as follows: define

$$S = \{(x_1, x_2) \in B : b_{1l}/a_{1l} \leq \|x_1\| \leq b_1^u/a_1^u,\ b_{2l}/a_{2l} \leq \|x_2\| \leq b_2^u/a_2^u,\}, \quad \textbf{(3.1)}$$

And let

$$K^* = \{\lambda(x_1, x_2) \in B : (x_1, x_2) \in S, \lambda \geq 0\}. \quad \textbf{(3.2)}$$

Since $\hat{G}_2/G_1 \subset K^*$, it is not difficult to see that N: $\hat{G}_2/G_1 \to B$ has at least one fixed point which is a periodic solution of (2.1)-(2.4). We will prove that this solution is unique by applying the result of Smith cited above. The following lemmas are needed first.

LEMMA 4. If $(x_1, x_2) \in K^* - \{0.0\}$ and if $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

PROFF. If $(x_1, x_2) = N(x_1, x_2)$ and $(x_1, x_2) \in K^* - \{0,0\}$, then $(x_1, x_2)$ is a periodic solution of (2.1) –(2.4) and hence we have

$$(b_{1l} - a_{1l}x_1)x_1 < \frac{dx_1}{dt} < (b_1^u - a_1^u x_1)x_1, \qquad (\mathbf{3.3})$$
$$(b_{2l} - a_{2l}x_2)x_2 < \frac{dx_2}{dt} < (b_2^u - a_2^u x_2)x_2,$$

And hence any nonzero periodic solution $(x_1, x_2)$ of (2.1)-(2.4) has to be such that

$$b_{1l}/a_{1l} < x_1(t) < b_1^u - a_1^u,$$

$$b_{2l}/a_{2l} < x_2(t) < b_2^u - a_2^u,$$

For $t > -\tau$, which proves the lemma.

LEMMA 5. The mapping $N$ defined by (2.5)-(2.9) is monotone nondecreasing on $K^*$.

PROOF. Let $(x_1, x_2) \in K^*$, $(y_1, y_2) \in K^*$ such that $(x_1 - y_1, x_2 - y_2) \in K^*$ and $(x_1, x_2) \neq (y_1, y_2)$. Then $x_1 \geq y_1$, $x_2 \geq y_2$. From the definition of $N$ we derive that

$N_1(t, x_1, x_2) - N_1(t, y_1, y_2)$

$$= \int_0^{\omega} G_{11}(m, \omega, t-s)\{m[x_1(s)] + x_1(s)f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s)f_1(s, y_1(s), y_2(s-\tau_{12}))\}ds$$

$$+ \int_0^t G_{12}(m, 0, t-s)$$

$$x\{m[x_1(S) - y_1(s)] + [x_1(s)f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s)f_1(s, y_1(s), y_2(s-\tau_{12}))]\}ds \qquad \mathbf{(3.4)}$$

With a similar relation for $N_2(t, x_1, x_2) - N_2(t, y_1, y_2)$. Since $f_1$ is continuous, $f$ is bounded on $K^*$; also it will follow from $(x_1, x_2) \neq (y_1, y_2)$ that

$$0 \not\equiv x_1(s) - y_1(s) \geq 0 \text{ and } 0 \not\equiv x_2(s) - y_2(s) \geq 0;$$

A sufficiently large positive constant $m$ will make the right side of (3.4) positive.A similar argument will lead to $N_2(t, x_1, x_2) - N_2(t, y_1, y_2) > 0$ perhaps with a larger constant $m$. Finally we can choose the bigger of these two constants $m$ to yield

$$N(x_1, x_2) - N(y_1, y_2) \text{ for } (x_1 - y_1, x_2 - y_2) \in K^* - \{(0,0)\}.$$

The monotonicity of $N$ on $K^*$ follows.

LEMMA 6. Let $(x_1, x_2) \in$ interior of $K^*$ and let $\lambda$ be $a$ constant such that $0 < \lambda < 1$. Suppose $f_1, f_2$ satisfy (2.2) - (2.4). Then there exists $a$ Positive constant $\eta$ such that

$$N(\lambda x_1, \lambda x_2) \geq \lambda(1+\eta) N(x_1, x_2). \quad \textbf{(3.5)}$$

PROOF. A direct computation leads to

$$N_1(t, \lambda x_1, \lambda x_2) - \lambda N_1(t, x_1, x_2)$$

$$= \lambda \int_0^{\omega} G_{11}(m, \omega, t-s) g_1\big(s, x_1(s), x_2(s-\tau_{12})\big) H_1(s, \lambda, x_1, x_2) ds$$

$$+ \lambda \int_0^{t} G_{11}(m, 0, t-s) g_1\big(s, x_1(s), x_2(s-\tau_{12})\big) H_1(s, \lambda, x_1, x_2) ds$$

**(3.6)**

Where

$$H_1(s, \lambda, x_1, x_2) = \frac{m\, x_1(s) + x_1(s) f_1\big(s, \lambda\, x_1(s), \lambda\, x_2(s-\tau_{12})\big)}{m\, x_1(s) + x_1(s) f_1\big(s, x_1(s), x_2(s-\tau_{12})\big)} - 1. \quad (\mathbf{3.7})$$

Since $f_1$ is decreasing (see (2.2)) with respect to $x_1$ and $x_2$ there exists a positive constant say $\eta_1$ such that

$$H_1(s, \lambda, x_1, x_2) \geq \eta_1 \text{ for } 0 < \lambda < 1, s > 0, (x_1, x_2) \in K^*, \quad \textbf{(3.8)}$$

And hence

$N_1(t, \lambda x_1, \lambda x_2) \geq \lambda(1+\eta_1)N_1(t, x_1, x_2)$.

A similar argument will show that there exists an $\eta_2 > 0$ such that

$N_2(t, \lambda x_1, \lambda x_2) \geq \lambda(1+\eta_2)N_2(t, x_1, x_2)$,

And the result follows if we choose $\eta = \min(\eta_1, \eta_2)$.

Now the following uniqueness result is a direct consequence of the result of Smith [11] and the Lemmas 4, 5, 6.

THEOREM 2. *Suppose $f_1, f_2$ of (2.1) satisfy (2.2) – (2.4). Then there exists a unique $\omega$ - periodic strictly positive periodic solution of (2.1).*

## 4. Linear stability of the periodic solution

Let $\{x_1^*(t), x_2^*(t), t \geq -\tau\}$ be the strictly positive unique periodic solution of (2.1). Let $\{x_1(t), x_2(t), t \geq -\tau\}$ be any other solution of (2.1). To study the linear stability of $(x_1^*, x_2^*)$ we let

$$\begin{aligned} x_1(t) &\equiv x_1^*(t)\{1 + X_1(t)\}, \\ x_2(t) &\equiv x_2^*(t)\{1 + X_2(t)\}, \end{aligned} \quad t \geq -\tau, \qquad \textbf{(4.1)}$$

And derive that the perturbations $(x_1, x_2)$ are governed by the linear variational system (after neglecting the non linear terms in $x_1, x_2$)

$$\begin{aligned} \frac{dX_1(t)}{dt} &= -a_{11}(t)\, x_1^*(t)X_1(t) - a_{12}(t)x_2^*(t-\tau_{12})\, X_2(t-\tau_{12}) \\ \frac{dX_2(t)}{dt} &= -a_{21}(t)x_1^*(t-\tau_{21})x_1(t-\tau_{21}) - a_{22}(t)x_2^*(t)x_2(t), \end{aligned} \qquad \textbf{(4.2)}$$

For $t > 0$, where

$$\left.\begin{aligned} -a_{11}(t) &= \frac{\partial f_1}{\partial x_1} \\ -a_{12}(t) &= \frac{\partial f_1}{\partial x_2} \end{aligned}\right\} \text{Evaluated at } (t, x_1^*(t), x_2^*(t-\tau_{12})),$$

**(4.3)**

$$\left.\begin{aligned} -a_{22}(t) &= \frac{\partial f_2}{\partial x_1} \\ -a_{22}(t) &= \frac{\partial f_2}{\partial x_2} \end{aligned}\right\}$$ Evaluated at $(t, x_1^*(t-\tau_{21}), x_2^*(t))$.

To establish the linear asymptotic stability of $(x_1^*(t), x_2^*(t))$ it is enough to prove that of the trivial solution of (4.2) - (4.3) We will do this by means of a combination of a comparison differential system and a Lyapunov function. First we define

$$\frac{d|X_1(t)|}{dt} = \begin{cases} \frac{dX_1(t)}{dt} & if\ X_i(t) \geq 0, \\ -\frac{d\setminus X_1(t)}{dt} & if\ X_i(t) < 0, \end{cases} \quad i = 1,2, \qquad \textbf{(4.4)}$$

And derive from (4.2)- (4.3) the following comparison system of differential inequalities:

$$\begin{aligned} \frac{d|X_1(t)|}{dt} &\leq -b_{11}, x_1^*(t)\,|X_1(t)| + b_{12}x_2^*(t-\tau_{12})\,|X_2(t-\tau_{12})|, \\ \frac{d|X_2(t)|}{dt} &\leq -b_{21}, x_1^*(t-\tau_{21})\,|X_1(t-\tau_{12})| - b_{22}x_2^*(t)\,|X_2(t)|, \end{aligned} \qquad \textbf{(4.5)}$$

For $> 0$, where

$$-b_{11} = \max_{t\geq -\tau}\{-a_{11}(t)\},$$

$$\begin{aligned} b_{12} &= |\min_{t\geq -\tau}\{-a_{12}(t)\}| \\ b_{21} &= |\min_{t\geq -\tau}\{-a_{21}(t)\}| \end{aligned} \qquad \textbf{(4.6)}$$

$$-b_{22} = \max_{t\geq -\tau}\{-a_{22}(t)\}$$

We can now formulate the result as follows. By the periodicity of $a_{ij}$, the constants $b_{ij}$ are independent of $\tau$ (see also (2.2)).

THEOREM 3. *The unique strictly positive periodic solution* $\{x_1^*(t), x_2^*(t), t \geq -\tau\}$ of $(2.1) - (2.4)$ *is linearly asymptotically stable provided*

$$|-b_{11}| > b_{21} \text{ and } |-b_{22}| > b_{12},$$

**(4.7)**

*Where the constants* $b_{ij}$ $(\mathrm{i},\mathrm{j} = 1,2)$ *are as defined in* (4.6).

PROOF. Consider the scalar valued Lyapunov-like function $v(t) = V(t, X_{1,}X_{2,})$where

$$v(t) = V(t, X_{1,}X_{2,}) = |X_{1,}(t)| + |X_{2,}(t)| \int_{t-\tau_{12}}^{t} b_{12} x_2^*(u)|X_{2,}(u)|du$$

$$+\int_{t-\tau_{21}}^{t} b_{21} x_1^*(u)|X_{1,}(u)|\mathrm{du}. \qquad \textbf{(4.8)}$$

Since $x_1^*, x_2^*$ are strictly positive, it is found from (4.7)-(4.) that

$$\frac{dv}{dt} \leq -[(|b_{11}| - b_{21})\, x_1^*(t)|X_{1,}(t)| + (|b_{22}| - b_{22})\, x_2^*(t)|X_2(t)|] \qquad \textbf{(4.9)}$$

For $t > 0$. It follows from (4.7)-(4.9) that $V(t, X_{1,}X_{2,}) \leq |X_{1,}(t)| + X_{2,}(t)|$ and $dv/dt \leq 0$,which will imply the stability of the trivial solution of (4.2)-(4.3). We claim that (4.9) also implies the asymptotic stability of the trivial solution of (4.2)-(4.3)in the sense that

$$\lim_{t-\infty} [\,|X_{1,}(t)|+|X_{2,}(t)|\,] = 0,$$

**(4.10)**

For all bounded initial conditions $X_{1,}(s), X_{2,}(s), s \in \{\mathcal{T}, 0\}$. $Suppose (4.10) is\ not$ true Then we can find an increasing sequence $\{t_m\} \to \infty$ as $m \to \infty$ such that

$$(|b_{11}| - b_{21})\, x_1^*(t_m)|X_{1,}(t_m)| + (|b_{22}| - b_{12} x_2^*(t_m)|X_2(t_m)| > \varepsilon,$$

M=1, 2, 3…

**(4.11)**

Using the right derivative *dv/dt* of $v$ we derive that

$$v(t_m - \varepsilon) - \mathrm{v}(t_m) \leq -\int_{t_m-\varepsilon}^{t_m} \frac{\mathrm{dv}(\mathrm{T})}{\mathrm{dT}}\,\mathrm{d}\tau.$$

**(4.12)**

If $\varepsilon$ is small enough we have from (4.9)(4.11),(4.12)that $dv(\tau)/dr < (-\varepsilon/2)$ and

$$v(t_m - \varepsilon) - \mathrm{v}(t_m) \geq \varepsilon^2/2 \text{ and hence}$$

$$v(t_m) \geq v(t_m - \varepsilon) - (\varepsilon^2/2)$$

$$\geq v(t_m - 1)\text{-}(\varepsilon^2/2)$$

$$\geq v(t_m - 2)\text{-}(\varepsilon^2/2)$$

$$\geq v(t_0)\text{-m}\,(\varepsilon^2/2)$$

Thus for sufficiently large $m, v(t_m)$becomes negative which contradicts the Nonnegative nature of $v\ for\ t \geq 0$.

this contradiction shows that $|X_{1,}(\mathrm{t})| + |X_2(\mathrm{t})| \to 0$ as $\mathrm{t} \to \infty$ and the proof is complete.

As an example let us consider the case of two species Lotka- Volterra type competition in a periodic environment for for which we have

$$f_1\big(\mathrm{t}, x_1(\mathrm{t}), x_2(\mathrm{t} - T_{12})\big)1 \quad \left\{ \frac{x_1(\mathrm{t}) + ax_2(\mathrm{t} - T_{12})}{K_1(\mathrm{t})} \right\}$$

$$f_2(\mathrm{t}, x_1(\mathrm{t} - {}_{21})x_2(\mathrm{t})) = 1 - \left\{ \frac{\beta x_1(\mathrm{t}-T_{21}) + x_2(\mathrm{t})}{K_2(\mathrm{t})} \right\}$$

Where $\alpha, \beta$ are positive constants and $k_{1,}k_{2,}$ are strictly positive continuous periodic functions with a common period Equation (4.6)and (4.7) for this special case lead to

$$\left|\operatorname{axm}_{t>T}\{-a_{11}(t)\}\right| = \max_{t>T}(1/K_1(t)) > \min_{t>T}\left(-a_{21}(\mathrm{t})\right)\Big| = \min_{t\in}, \beta/K_2(\mathrm{t})$$

$$\max_{t\in T}\{-a_{22}(t)\}\Big| = \max_{t\in T}(1/K_2(t)) > \min_{t\in T}\left(-a_{12}(\mathrm{t})\right)\Big| = \min_{t\in} \mathfrak{a}/K_1(\mathrm{t})$$

The inequalities reduce in the absence of temporal variations in $k_{1,}k_{2,}$ to the inequality 0< $\mathfrak{a}$,$\beta$<1, which is well known for the global asymptotic stability of the positive steady state corresponding to the above example in a constant systemic domain the above example shows also that delays in interspecific interactions have no influence on the local asymptotic stability of the unique periodic solution of the periodic ecosystem A number of cases of this type in a constant environment have been established recently by the author (Gopalsamy [6]).A numerical example is illustrated at the end.

**Consequences**

It can be found that all the conclusions on the existence uniqueness and stability of a periodic solution of (2.1)-(2.4) are valid when $OT_{12}$=0 or $T_{12}$=0 or both $T_{12}$=0 and $T_{21}$=0: thus it will follow that delays in interspecific competition do not destabilize an otherwise stable coexistence (oscillatory or nonoscillatory) of two competing types of interacting particles atoms molecules transactions interaction. In forthcoming work we will extend the analysis of the system (2.1)-(2.4) to a system of n-competing species with a hypothesis similar to that of (2.4) It is however an open problem to investigate the effects of delays in intraspecific competition in a periodic milieu

Pianka ([8],pages 139-140) remarks that “In real population rate of increase competitive abilities and carrying capacities do very from individual to individual with population density in space and time Indeed temporal variation in the environment may often allow coexistence by continually altering the competitive abilities of population inhabiting it Moreover time lags are doubt lessly of some importance in real population Finally a heterogeneous environment may allow real competitors to evolve divergent resource utilization patterns and to reduce interspecific competitive inhibition

To understand the mechanism of selection pressure against competitive inhibition we first write the Malthusian parameters $\mu_1$, $\mu_2$ of fitness for the two species as follows (see MacArthur[7],page 226-230):

$$\mu_1=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_1\left(t,x_1(t),x_2(t-T_{12})\right)$$

$$\mu_2=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_2(t,x_1(t-T_{21}),x_2(t))$$

Since $\left(x_1(t),x_2(t)\right)\rightarrow\left(x_1^*(t),x_2^*(t)\right)$ for large t we can write the parameters of finess in the from

$$\mu_1=f_1\left(t,x_1^*(t),x_2^*(t-T_{12})\right)$$

$$\mu_2=f_2\left(t,x_1^*(t-T_{12}),x_2^*(t)\right)$$

In these relations we can use

$$x_2^*(t-T_{12}),x_2^*(t)-T_{12}\frac{dx_2^*(t-T_{12}+\theta_1T_{12}}{dt}$$

$$x_1^*(t-T_{12}),x_1^*(t)-T_{21}\frac{dx_1^*(t-T_{21}+\theta_2T_{21}}{dt}$$

$\theta_1,\theta_2,\in(0,1)$ and rewrite $\mu_1,\mu_2$ in the form

$$\mu_1=f_1\left(t,x_1^*(T),x_2^*(t)-T_{12}\frac{dx_2^*(t-T_{12}+\theta_1T_{12}}{dt}\right.$$

$$\mu_2=f_2\left(t,x_1^*(T),-T_{21}\frac{dx_1^*(t-T_{21}+\theta_2T_{21}}{dt},x_2^*(t)\right)$$

$\theta_1,\theta_2,\in(0,1)$

It is seen that the presence of the delays $T_{12}$ and $T_{21}$ provides additional sources of variability in the finess parameters and thereby in the selection pressure against interspecific competition Since a necessary and sufficient condition for natural selection to act is differential (variable) rates of confederation growth we find that time delays coupled with periodicity of an domain may allow real to evolve divergent resource utilization patterns and to reduce

Interspecific competitive inhibition for some remarks relating to the significance of a periodic environment to sustain the coexistence of a large number of confederation types in the absence of delays we refer to a recent article (Gopalsamy)The following example has been numerically solved and the solutions are displayed graphically for two different values of the delay parameter T:

$\frac{dx_1(t)}{dt} x_1(t)[1 - x_1(t) + 0.4x_2(t - \mathcal{T})/k_1(t)]$

$\frac{dx_2(t)}{dt} x_2(t)[1 - \{0.8\, x_1(t - \mathcal{T}) + x_2(t)\}/k_2(t)]$

Where $k_1(t) = 20 + \sin(\pi t)$ and $k_2(t)$=2.5+sin (πt)

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There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. “All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ....” (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies

systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational

content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification.

## GOPALSAMY'S MODEL APPIED TO COSMOLOGICAL PROBLEMS

## GLOBAL ASYMPTOTIC STABILITY IN AN ALMOST PERIODIC LOTKA- VOLTERRA SYSTEM

### Abstract

**Gopalsamy studied** Sufficient conditions are obtained for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of a Lotka- Volterra system with almost periodic coefficients. Research was published Australian. Math Soc. B.27 (1986) 346-360. Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of

decay rate of perturbations. Sole intention, main objective and primary concern of the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are counted in the same vein as that of predator prey. Note that the model is applicable to each and every module. Quintessentially each parameter should have been given different name such as superscript 1,2, etc., as that has been done in the foregoing and a formal concatenation done at the end. Due to spatial restrictions we have not been able to do the line of thought and manner of analysis as has been done hereinbefore. This is the most important aspect of the model. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 11. Introduction

At the outset we assume that the system. This article is concerned with the derivation of a set of "easily verifiable" sufficient conditions for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of the Lotka- Volterra system

$$\frac{dx_1(t)}{dt} = x_i(t)\left\{b_1(t) - \sum_{j=1}^{n} a_{ij}\,(t)x_j(t)\right\};$$

$$t>t_0;\ t_0\ \in (-\infty,\infty);\ i = 1,2,\dots,n, \qquad (1.1)$$

where $b_1, a_{ij}\ (i,j = 1,2,\dots.,n)$ are nonnegative almost periodic functions defined for $t \in (-\infty,\infty)$.

In question satisfies Haimovici's equation based on age wise stratification given in the foregoing. Note that Haimovici's equations can be applied to Lotka-Volterra system term by term.

Let R and $R_n$ respectively denote the set of all real numbers and the $n$ - dimensional real Euclidean space; $R_n^+$ will denote the nonnegative cone of $R_n$. Since almost- periodic functions defined on R are bounded, we define the constants $b_i^l, b_l^u, a_{ij}^l, a_{ij}^u, (i,j = 1,2,\dots,n)$ by the following:

$$\inf_{t\in R} b_l(t) = b_i^l; \qquad \sup_{t\in R} b_l(t) = b_l^u, \qquad i,j = 1,2,\dots,n.$$

$$\inf_{t\in R} a_{lj}(t) = a_{lj}^l; \qquad \sup_{t\in R} a_{lj}(t) = a_{ij}^u.$$

We shall study the almost- periodic system (1.1) with the following assumption on the coefficients of (1.1);

$$b_l^l > 0; \qquad a_{ll}^l > 0, \tag{1.2}$$

$$b_l^l > \sum_{\substack{j=1 \\ j\neq i}}^{n} a_{lJ}^u (b_j^u / a_{iJ}^l). \tag{1.3}$$

Since the solutions of (1.1) corresponding to nonnegative initial conditions remain nonnegative subsequently, it will follow that

$$\frac{dx_1}{dt} \leq x_l\{b_l^u - a_{ll}^l x_i\}; \qquad t > t_0, i = 1,2,\dots,n, \tag{1.4}$$

As a consequence of which we will have

$$0 < x_l(t_0) \leq b_l^u / a_{ll}^l = x_i^u \Rightarrow x_i(t) \leq b_l^u / a_{ll}^l, \qquad i = 1,2,\dots,n; t \geq t_0. \tag{1.5}$$

(1.6) And (1.5) together lead to

$$\frac{dx_1}{dt} \leq x_l \left\{ b_l^u \sum_{\substack{j=1 \\ j\neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right) - a_{ll}^l x_i \right\}, \qquad i = 1,2,\dots,n; \tag{1.6}$$

Now (1.2), (1.3) and (1.6) lead to

$$x_i(t_0 \leq \left( b_i^l - \sum_{\substack{j=1 \\ J\neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right) \right) (1/a_{ll}^u) = x_i^l$$

$$\Rightarrow x_l(t) \to x_i^l \ for\ t \to t_0,\ i = 1, 2, \dots, n. \quad (1.7)$$

12. It is our aim to show that the mode together with other models stated in the foregoing together with haimovicis model is applicable to whole lot of consummative debit – credit systems in cosmology, astrophysics, finance, economics, sociology, anatomy computation. We will not however divide the system studied here into 9 components/ modules due to restrictions on space….

From the foregoing preparation we have the following:

LEMMA 1.1. *Assume that* $b_i, a_{ij} (i, j = 1, 2, \dots, n)$ *are scalar nonnegative almost periodic functions defined on* $\boldsymbol{R}$ *such that (1.2) and (1.3) hold. Then the set S defined by*

$$S = \{x = (x_1, x_2, \dots, x_n) \in \boldsymbol{R}_n | x_* \to x_l \to x^*; i = 1, 2, \dots, n\} \quad (1.8)$$

*Where*

$$x_* = \min_{1 \to n} x_l^l; \qquad x_* = \min_{1 \to n} x_i^u \quad (1.9)$$

*is invariant with respect to (1.1)*

**Appellation, cognomen, compellation, designation, denomination of the conducive, infelicitous and unfortutious** responses to environment change may be influenced by changes in available domain and milieu under consideration. An example is the following study: Authors describe the gravitational degrees of freedom of the Schwarzschild black hole by one free variable. Authors introduce an equation which we suggest to be the Schrödinger equation of the Schwarzschild black hole corresponding to this model. We solve the Schrödinger equation explicitly and obtain the mass spectrum of the black hole as such as it can be observed by an observer very far away and at rest relative to the black hole. Our equation implies that there is no singularity inside the Schwarzschild black hole, and that the black hole has a certain ground state in which its mass is non-zero. **Schrodinger equation of the Schwarzschild black hole Jarmo Mäkelä (Cambridge U E-Print: gr-qc/9602008 | PDF** Particles, axions, atoms, molecules, quarks, Majorana fermions, bosons et al in its confederation capacity and aggregation processes, type interactions and interactions between landscape dynamics. Following bears ample testimony, infallible observatory to the fact there might be changes in the quantum behaviour due to merger of black holes. Following example bears enough testimony, infallible observatory and impeccable demonstration and apodictic evidence as how some of the models could be applied to cosmological problems. This is attributable and ascribable to the fact that the classification criterion is based on the age.

**How gravity kills (e) Schrödinger's cat Theorists argue that warped space-time prevents (e) quantum superpositions of large-scale objects. Elizabeth Gibney 17 June 2015 Propositions are written in model able form and can easy be scheduled.**

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| (111) If the cat in Erwin Schrödinger's famous thought-experiment behaved according to quantum theory, it would be able to exist in multiple states at once: both dead and alive. |
| (112) Physicists' common explanation for why we don't (e) see such quantum superpositions — in cats or any other aspect of the everyday world — is interference from (e) the environment. |
| (113) As soon as a quantum object interacts with (e&eb) a stray particle or a passing field, it picks (e) just one state, collapsing (eb) into our classical, everyday view. |
| (114) But even if physicists could completely isolate (eb) a large object in (eb) a quantum superposition, according to researchers at the University of Vienna, it would still collapse into (e&eb) one state — on Earth's surface, at least. |
| (115) "Somewhere in interstellar space it could be that the cat has (e) a chance to preserve quantum coherence, but on Earth, or near any planet, there's (=) little hope of that," says Igor Pikovski. The reason, he asserts, is gravity. |
| (116) Pikovski and his colleagues' idea, laid out in a paper published in Nature Physics on 15 June1, is at present only a mathematical argument. But experimenters hope to test whether gravity really does collapse (eb) quantum superpositions, says Hendrik Ulbricht, an experimental physicist at the University of Southampton, UK. "This is a cool, new idea, and I'm up for trying to see it in experiments," he says. Assembling the technology to do so, however, may take as long as a decade, he says. |
| (117) Cinema-goers who saw the film Interstellar are already familiar with the basic principle behind the Vienna team's work. Einstein's theory of general relativity states that an extremely massive object causes clocks near it to (e) run more slowly because (e) its strong gravitational field stretches (e&eb) the fabric of space-time (which is why a character in the film aged only an hour near a black hole, while (e) seven years passed on Earth). On a subtler scale, a molecule placed nearer the Earth's surface experiences (e&eb) a slightly slower clock than one placed slightly further away. |
| (118) Because of gravity's effect on (e&eb) space-time, Pikovski's team realised that variance in a molecule's position will also influence (e&eb) its internal energy — the vibrations of particles within (eb) the molecule, which evolve (eb) over time. |
| (119) If a molecule were put in a quantum superposition of two places, the correlation between position and (e&eb) internal energy would soon cause (eb) the duality to 'decohere' to (e) the molecule taking just one path, they suggest. |
| (120) "In most situations decoherence is due to (e) something external; here it's as though the internal jiggling is interacting with (e&eb) the motion of the molecule itself," adds Pikovski. |

**A practical limit**

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| (121) No one has yet seen this effect because other sources of decoherence — such as magnetic fields, thermal radiation and vibrations — are (=) typically much stronger and cause (eb) quantum systems to collapse long before gravity becomes an issue. But experimenters are keen to try. |
| (122) Markus Arndt, an experimental physicist also at the University of Vienna, already tests whether quantum superpositions can be observed for (e) large, though not cat-sized, objects (see Fattening up Schrödinger's cat). |
| (123) He sends large molecules through (e&eb) a matter-wave interferometer, a system that gives (eb) each molecule a choice of (e) two different paths. |
| (124) In the classical view, a molecule travels down just one path; a quantum molecule effectively passes through (e&eb) both routes at once and interferes with (e&eb) itself to create (eb) a characteristic wave-like pattern (see image, below). |

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| (125) | A quantum interference pattern built up from complex molecules known as phthalocyanines.(deleted due to spatial constraints) |
| (126) | A similar set up could be used to (e) test for gravity's ability to destroy (e) quantum behaviour: by comparing a vertical interferometer, in which the superposition should soon decohere due to (e) the stretching of time in one path versus (e&eb) the other, with a horizontal set up, where (e) the superposition could remain. |
| (127) | Arndt, who has tested the effect for (e) molecules of up to 810 atoms2, points out those large molecules, would be (=) good for testing the gravitational effect because (e) they contain (e) lots of particles contributing to (e) internal energy. |
| (128) | **But not only would the researchers have to further suppress the external environment to reduce (e) other decoherence effects, they would also need to (e) increase the separation of the two paths from micrometres to metres, or else use molecules perhaps 1 million times more massive. "It's certainly very challenging," says Arndt**. |
| (129) | If gravity's effect does limit (e) quantum behaviour on Earth, tests of quantum reality for (e) large objects may eventually have to (e) move into space, says Angelo Bassi, a physicist at the University of Trieste in Italy. |
| (130) | "But from a deep, fundamental point of view, this is nothing new," he says. A gravitational field is merely another environment to interact with, so invoking it does not (e) explain whether quantum behaviour might lead to (eb) classical reality if gravity's influence were mitigated — for example, by doing the experiment in (eb) gravity-free space. |
| (131) | Nor does the effect described by Pikovski and colleagues say anything about quantum gravity: a theory that would unite gravity and (e&eb) quantum mechanics in (eb) a single description, which many researchers are working towards. |
| (132) | "It's an interesting effect, but it's still quantum physics applied to (e&eb) classical general relativity. In that way it doesn't change our picture of the world," adds Bassi. Nature doi:10.1038/nature.2015.17773 |

Integrated factors these influences because they deal with ambience, mise en scene change or confederation dynamics, but rarely both. Time series of (e) space time suitability models with spatially explicit stochastic models to explore factors that influence the viability of type viability and confederation feasibility under stable and changing scenarios. Methodology could be applied to interrogation propositions and problems which might subsume one with the other notwithstanding that there is a space of singularities and there is space for completion of present from the past and future. An example is Time image movement from Deleuze. We clearly face a stiff challenge in our attempt to proffer Deleuzian treatments of photography as our critical project brushes up against the fact that Deleuze does not invest photography with (e&eb)the creative and re-creative potential that he associates with (e&eb) both cinema and modern painting. In short, while he highly values other visual artistic forms, he seemingly presents photographic texts as stagnate documents or tools that produce certainty, organize bodies and desires, and iterate hackneyed ideas. He even uses photography as something of a foil to demonstrate the innovation of the cinema and the originality of modern painters. In Francis Bacon: The Logic of Sensation (1981), he notes how "photography has taken over the illustrative and documentary role, so that modern painting no longer needs to fulfill this function, which still burdened earlier painters" (10). He later explains that "photographs are ways of seeing, and as such, they are illustrative and narrative reproductions or representations ... . They are what are seen, until finally one sees nothing else." He treats the photo as an instrument for reproducing representations of reality—a device that iterates images until they are ossified as established stories, icons, or even stagnant perceptions. Indeed, like Bergson before him, Deleuze compares photography to what he identifies as a reductionist model of perception: "what we see, what we perceive, are photographs" (74).[4] The photo is, in effect, an always-already passé

sensational experience that numbs our sensitivity to the ongoing vitality of life and provides us with an efficient understanding of the dynamism of our world. But it is compelling that Deleuze theorizes both the quotidian nature and the liberatory potential of photography; he treats it as a common, everyday visual experience that, despite its aesthetic and technological limitations, emancipates the modern painter to explore new artistic opportunities. This special issue of Rhizomes repeatedly provides artistic examples and critical strategies that invite us to see photography as replete with the power to unburden and emancipate that Deleuze associates with both film and painting. Indeed, one way to think of this project of reading Deleuze against Deleuze might be to ask: if photography can free painting and painters to pursue novel creative opportunities, can it likewise free itself, photographic practitioners, and even viewers to create and re-create anew? **Rhizomes » Issue 23 (2012) » Michael Kramp Unburdening Life, or the Deleuzian Potential of Photography Michael Kramp Lehigh University** dmk209@lehigh.edu

**Existence of an almost- periodic solution**

DEFINITION 2.1. A function $f(t,x)$,where $f$ is an $m$- vector, $t$ is a real scalar and $x$ is an $n$- vector, is said to be almost periodic in $t$ uniformly with respect to $x \in X \subset R_n$, $if\ f(t,x)$ is continuous in $t \in R\ and\ x \in X$, and if for any $\varepsilon > 0$, it is possible to find a constant $l(\varepsilon) > 0$ such that in any interval of length $l(\varepsilon)$ there exists a $\tau$ such that the inequality

$$\left||f\ (t+\ \tau,x) - \ f(t,x)|\right| = \sum_{l=1}^{m} |f_l\ (t+\ \tau,x) = \ f_l(t,x)|$$

$$< \varepsilon \qquad (2.1)$$

is satisfied for all $t \in (-\infty,\infty), x \in X$. The number $\tau$ is called an $\varepsilon -$ translation number of $f(t,x)$.

Let $\{\lambda_j\}$ denote the set of all real numbers such that

$$\lim_{T\to\infty} T^{-1} \int_0^T f(t,x)\exp(-i\,\lambda_j t)dt \not\equiv 0 \quad for\ x \in X; i = \sqrt{-1}. \qquad (2.2)$$

It is known that when the set X is separable, the set of numbers $\{\lambda_j\}$ in (2.2) is countable. The set $\{\sum_1^N n_j\ \lambda_j\}$ for all integers N and integers $n_j$ is called the module of $f(t,x)$.

The following result concerned with module containment will be used in the proof of our existence theorem below.

**MODULE CONTAINMENT THEOREM** (Yoshizawa [7], page 18). *Let $f(t,x)$ and $g(t,x)$ be almost periodic in t uniformly for $x \in D \subset \boldsymbol{R}_n$. If for any compact set $S \subset D$ and for any sequence of real numbers $\{\tau_k\}$ having a limit as $k \to \infty$ (the limit being finite or$\pm\infty$) for which $\{f(t+\tau_k,x)\}$ is uniformly convergent on $\boldsymbol{R} \times$*

$S, \{g(t+,x)\}$ *is also uniformly convergent on* $\boldsymbol{R} \times S$ *then the module of* $g(t+,x)$ *is contained in the module of* $f(t+,x)$.

We can now formulate our existence theorem.

**THEORM 2.1** *Assume that the almost-periodic parameters* $b_1, a_{1J}$ $(i,j = 1,2,\dots,n)$ *of (1.1) satisfy (1.2)-(1.3). Furthermore assume that there exists a positive number* $\mu$ *such that*

$$\inf_{1 \in \boldsymbol{R}} a_{11}(t) > \sum_{\substack{J=1 \\ J \neq 1}}^{n} (\sup_{t \in \boldsymbol{R}} a_{Jl}(t)) + \mu, \quad i = 1,2,\dots,n. \tag{2.3}$$

*Then the system (1.1) has a strictly positive (component wise) almost-periodic solution, say* $w(t) = w_1(t), \dots, w_n(t), t \in \boldsymbol{R}$, *whose module is contained in that of*

$$f(t,x) = \{f_1(t,x_1,\dots,x_n), f_2(t,x_1,\dots,x_n), \dots, f_n(t,x_1,\dots,x_n)\}$$

*Where*

$$f_l(t,x_1,\dots,x_n) = x_l \left\{ b_l(t) - \sum_{j=1}^{n} a_{lj}(t) x_J \right\},$$

$$i = 1,2,\dots n; (x_1,x_2,\dots,x_n) \in S. \tag{2.4}$$

PROOF. The set $S$ defined by (1.8)-(1.9) is compact in $\boldsymbol{R}_n$. Let $\{\tau_h\}$ be any sequence of real numbers such that $\tau_k \to \infty$ as $k \to \infty$ and

$$\sum_{l=1}^{n} |f_i(t+\tau_k, x_1, \dots, x_n) - f_l(t, x_1, \dots, x_n)| \to 0 \quad \text{as } k \to \infty \tag{2.5}$$

Uniformly in $(x_1, \dots, x_n) \in S$ for $t \in (-\infty, \infty)$. for any real number $\beta$let $k_0(\beta)$ be the smallest value of $k$ such that $\tau k_0 + \beta \geq t_0$. Since the invariance of $S$ implies that any solution of (1.1)

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), x_2(t+\tau_k, t_0, x_2(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\} \in S$$

For all $t \geq \beta$ and $k \geq k_0(\beta)$ whenever $\{x_1(t_0), \dots, n_n(t_0)\} \in S$, our first task in the proof is to show that the sequence

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Converges to a continuous bounded function say $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined for $t \in [\beta, \infty)$ such that $w_1(t) \in [x_*, x^*], t \in [\beta, \infty), i = 1, 2, \dots, n$ and the convergence is uniform on all compact subsets of $[\beta, \infty)$. Since $x_*$ and $x^*$ are independence of $\beta$, it is enough to show that

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Is a Cauchy sequence on compact subsets of $[\beta, \infty)$.

Let $U$ be any compact subset of $\beta, \infty)$., and let be any arbitrary positive number. Choose an integer $n_0(\varepsilon, \beta) \geq k_0(\beta)$ so large that for $m \geq k \geq n_0(\varepsilon, \beta)$

and all $t \in (-\infty, \infty)$

$$\sum_{i=1}^{n} |b_1(t+\tau_m) - b_1(t+\tau_k)| < (\mu\varepsilon x_*)/(4x^*), \tag{2.6}$$

$$x^* \sum_{i=1}^{n}\sum_{j=1}^{n} |a_{lj}(t+\tau_m) - a_{lj}(t+\tau_k)| < (\mu\varepsilon x_*)/(4x^*) \tag{2.7}$$

$$(2nx^*/x_*)\exp[-\mu x_*(\beta + \tau_k - t_0)] < (\varepsilon/2x^*)$$

(2.8)

Denote $x_i(t, t_0, x_1, (t_0))$ $by$ $x_1(t)$ brevity, and consider a function $v(s)$ defined by

$$v(s) = v\left(x(s), x(s+\tau_m+\tau_k)\right)$$

$$= \sum_{i=1}^{n} |\log x_1, (s) - \log x_1(s+\tau_m+\tau_k)|, \qquad x \geq t_0, \tag{2.9}$$

Where $x(t) = \{x_1(t), \dots, x_n(t)\}$ denotes a solution of (1.1) with $\{x_1, (t_0), \dots, x_n(t)\} \in S$. the invariance of S for (1.1), together with the elementary mean value theorem of differential calculus, implies

$$(1/x^*)\sum_{l=1}^{n}|x_l(s)-x_1(s+\tau_m+\tau_k)|\ \upsilon(s)$$

$$(1/x^*)\sum_{i=1}^{n}|x_l(s)-x_1(s+\tau_m+\tau_k)|. \tag{2.10}$$

Calculating the right derivative $D^+\upsilon$ of $\upsilon$ along the solutions of (1.1) and simplifying

$$D^+\upsilon(s)=\sum_{i=1}^{n}\{sign[x_i(s)-x_i(s+\tau_m-\tau_k)]\}$$

$$\left\{\frac{d[logx_i(s)]}{ds}-\frac{d[logx_i(s+\tau_m-\tau_k)]}{ds}\right\}$$

$$\sum_{l=1}^{n}|b_l(s)-b_l(s+\tau_m-\tau_k)|$$

$$+\sum_{l=1}^{n}(-a_{il}(s)|x_i(s)-x_i(s+\tau_m-\tau_k)|+\sum_{\substack{J=1\\ j\neq l}}^{n}a_{ij}(s)|x_j(s)-x_j(s+\tau_m-\tau_k)|)$$

$$+\sum_{i=1}^{n}\sum_{j=1}^{n}|a_{ij}(s)|a_{ij}(s+\tau_m-\tau_k)|\,x_j(s+\tau_m-\tau_k) \tag{2.11}$$

$$-\mu x_*\upsilon\big(x(s),x(s+\tau_m-\tau_k)\big)+\sum_{l=1}^{n}|b_l(s)-b_i(s+\tau_m-\tau_k)|$$

$$+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|. \tag{2.12}$$

Let us set $s-\tau_k=\sigma$; then we have from (2.6)-(2.7) that

$$\sum_{i=1}^{n}|b_l(s)-b_1(s+\tau_m-\tau_k)|+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|.$$

$$= \sum_{l=1}^{n} |b_l(\sigma + \tau_k) - b_i(\sigma + \tau_m)|$$

$$+x^* \sum_{i=1}^{n} \sum_{J=1}^{n} |a_{ij}(\sigma + \tau_k) - a_{ij}(\sigma + \tau_m)|$$

$$< (\mu\varepsilon x^*)/(2x^*). \quad (2.13)$$

Now an integration of (2.12) over $(t_0, t - \tau_k)$ together with (2.13) leads to

$$v\big(x(t+\tau_k), x(t+\tau_m)\big) \leq v\big(x(t_0), x(t_0 + \tau_m - \tau_k)\big) \exp[-\mu x_*(t + t_k - t_0)]$$

$$+[(\mu\varepsilon x_*)/(2x^*)]/\mu x_*)$$

$$\leq [(2x_n^*)/x_*]\exp[-\mu x_*(t + t_k - t_0)] + (\varepsilon/2x^*)$$

$$< (\varepsilon/x^*) \text{ On using (2.8), (2.10)} \quad (2.14)$$

For $m \geq n_0(\varepsilon, \beta)$ and for all $t \in \cup \subset [\beta, \infty]$, Thus we have from the definition of $v$ and (2.14) that

$$\sum_{l=1}^{n} |log x_i(t + \tau_k) - log x_l(t + \tau_m)| < (\varepsilon/x^*) \quad (2.15)$$

And hence using $x_* \leq x_i(t) \leq x^*, i = 1, 2, \dots, n; t \geq t_0$,

$$\sum_{i=1}^{n} |x_i(t + \tau_k) - x_i(t + \tau_m)| < \varepsilon \text{ for all } t \in \cup \subset [\beta, \infty], m \geq n_0(\varepsilon, \beta). \quad (2.16)$$

The existence of a function $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined on $[\beta, \infty]$, such that $w_i(t) \in [x_*, x^*], i = 1, 2, \dots, n; t \in [\beta, \infty]$ follows from (2.16) and the fact that $x_i(t) \in [x_*, x^*], i = 1, 2, \dots, n;$ for $t \geq t_0$. Since $\beta$ is arbitrary, $w$ is in fact

Defined on $(-\infty, \infty)$ and hence we have

$x_i(t+\tau_k) \to w_l(t)$ as $k \to \infty$ uniformly on compact subsets of $(-\infty,\infty)$, $i = 1,2,\dots,n.$

(2.17)

Our next task is to show that $w: R \to S$ is differentiable and is a solution of (1.1). Since $x(t) = \{x_1(t),\dots,x_n(t)\}$ satisfies (1.1) we have $\sum_{i=1}^{n} |\frac{d}{dt} x_i(t+\tau_k) - \frac{d}{dt} x_i(t+\tau_m)|$

$$\sum_{i=1}^{n} |f_l\left(t+\tau_k, x_1(t+\tau_k),\dots, x_n\left(t+\tau_k\right)\right)$$

$$- f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)| \qquad (2.18)$$

$$+\sum_{l=1}^{n} |f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$- f_l\left(t+\tau_m, x_1(t+\tau_m),\dots,x_n(t+\tau_m)\right)|$$

Where $f_l\ (i = 1,2,\dots,n)$are as in (2.4). Since $x(t+\tau_k) = x_1(t+\tau_k),\dots, x_n(t+\tau_k) \in S$ for large $\tau_k$ for $t$ in each compact subset of $(-\infty,\infty)$ there exists an integer $n_1(\varepsilon) > 0$ such that if $m$ ? $k$ ? $n_1(\varepsilon)$, then for any $\varepsilon > 0$,

$$\sum_{i=1}^{n} |f_l\left(t+\tau_k, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$-f_1\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)| < (\varepsilon/2). \qquad (2.19)$$

Similarly since $x_l(t+\tau_k) \in S$ and $x_l(t+\tau_m) \in S$ for $i = 1,2,\dots,n$ there is an integer $n_2(\varepsilon) > 0$ such that for large enough $m, k$, such that $m$ ? $k$ ? $n_2(\varepsilon)$, we have

$$\sum_{l=1}^{n} |f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$-f_1\left(t+\tau_m, x(t+\tau_m),\dots,x_n(t+\tau_m)\right)| < (\varepsilon/2) \qquad (2.20)$$

This equation will be very much useful in chemistry if used with the modules mentioned and written in concatenation with each other. It should be borne in mind that the applicability is modulewise and in no case concatenation should be marginalized. In all cases parameters are different which could be represented with different superscripts and subscripts like 1,23,…….etc., **Anal Bioanal Chem. 2004 Sep; 380(1):36-45. Epub 2004 Jul 29. Electronic tongue for pharmaceutical analytics: quantification of tastes and masking effects Legin A et al**

The organoleptic aspects of pharmaceutical formulations affect their acceptability to the patient and hence can have an important effect on concordance with treatment. Objective evaluation of these aspects, particularly the taste of the formulation and the drug substance it contains, is difficult. Whilst volunteer taste panels can be used to good effect their utility is limited, particularly during very early stage development when the toxicological profile of the active pharmaceutical ingredient (API) is yet to be established in detail. A potentiometric "electronic tongue" has been applied to analyse a variety of 41 individual substances and mixtures of particular interest for pharmaceutical research and development. The electronic tongue (ET) was capable of discriminating between substances with different taste modalities and could also distinguish different substances eliciting the same basic taste; the ET is promising in terms of quantifying the content of each substance and has an ability to detect nuances of the basic taste (e.g. lingering or short-lived). After calibration the electronic tongue was successfully applied to predicting bitterness strength of binary mixtures with a sweetener in terms of "apparent" or "perceived" quinine content. In order to render a formulation palatable it is often necessary to mask the (usually bitter) taste of the API by the addition of masking agents such as sweeteners and flavours. The ET proved capable of distinguishing between formulations with different levels of sweetener and/or flavour in a manner that was consistent with their masking efficiency as perceived by a small human taste panel. A suitably calibrated ET could have the benefit of providing the pharmaceutical formulator with reliable data concerning the taste of the product quickly and with a reduced need to ask volunteers to taste active pharmaceutical samples. Early development activities could be facilitated when human tasting is usually not possible in the absence of the required toxicological data.

Thus if $m \neq n_0(\varepsilon) = max\{n_1, n_2\}$, we have

$$\sum_{l=1}^{n} |\frac{d}{dt}\, x_1(t+\ \tau_k) - \frac{d}{dt} x_1\ (\ (t+\tau_m)| < \varepsilon \qquad (2.21)$$

Which shows that

$$\lim_{k\to\infty} \frac{d}{dt} x_1(t+\ \tau_k) \quad exists\ for\ i = 1, 2, \dots, n$$

Uniformly on all compact subsets of $(-\infty, \infty)$. Thus we have

$$\lim_{k\to\infty} \lim_{k\to\infty} \frac{x_i(t+\tau_k+h) - \ x_i(t+\tau_k)}{h}$$

$$= \lim_{h\to 0} \lim_{k\to 0} \frac{x_l(t+\tau_k+h) - \ x_l(t+\tau_k)}{h} \qquad (2.22)$$

$$= \lim_{h\to 0} \frac{w_1(t+h) - \ w_1(t)}{h}, \qquad i = 1, 2, \dots, n,$$

Showing that (since the left side of (2.22) exists) $w_i (i = 1, 2, \dots, n)$ is differentiable. Furthermore we have from (2.22) that

$$\frac{dw_l(t)}{dt} = \lim_{k\to\infty} [\, f_i\ \big(t+\tau_k, x_1(t+\tau_k), \dots, x_n\ (t+\tau_k)\big)$$

$$-f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big)$$

$$+ - f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big)]$$

$$= \ -f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big); \qquad i = 1, 2, \dots, n, \qquad (2.23)$$

By the uniform almost periodicity of $f_1, f_2, \dots, f_n$ on $R \times S$ and hence $w(t) = \big(w_1(t), \dots, w_n(t)\big)$is a solution of (1.1).

In order to complete the proof, we have to show that $w(t) = \big(w_1(t), \dots, w_n(t)\big)$ is an almost- periodic function whose module is contained in that of $\{f_1(t, x_1, \dots, x_n), f_2(t, x_1, \dots, x_n), \dots, f_n(t, x_1, \dots, x_n)\}$ for $(x_1, x_2 \dots, x_n) \in S$. For this purpose it is in fact sufficient (by the module containment theorem) to show that for any sequence of real numbers $\{\tau_k\}$ for which

$$\{f_1(t + \tau_k, x_1, \dots, x_n), f_2(t, + \tau_k, x_1, \dots, x_n), \dots, f_n(t, + \tau_k, x_1, \dots, x_n)\}$$

Converges uniformly on R for all $(x_1, \dots, x_n) \in S$, the sequence $\{w_1(t, + \tau_k) \dots, w_n(t + \tau_k)\}$ converges uniformly on R where $\tau_k$ tends to a finite number or $\pm\infty$ as $k \to \infty$.

In the case where $\tau_k \to \tau \neq \pm\infty$ as $k \to \infty$, we have $w_l(t, + \tau_k) \to w_i(t + \tau), i = 1, 2, \dots, n$ uniformly on R as $k \to \infty$. Let us suppose $\tau_k \to \infty$ as $k \to \infty$. For any $\varepsilon > 0$ there exists an integer $n_0(\varepsilon)$ such that if $m$ ☒ $k$ ☒ $n_0(\varepsilon)$ then

$$\sum_{l=1}^{n} |b_l(t + \tau_k) - b_l(t + \tau_m)| < (\varepsilon/4x^*)x_*, \qquad (2.24)$$

$$x^* \sum_{i=1}^{n} \sum_{j=1}^{n} |a_{ij}(t + \tau_k) - a_{ij}(t + \tau_m)| < \left(\frac{\varepsilon}{4x^*}\right)\mu x_* \qquad (2.25)$$

$$n(2x^*/x_*) \exp[\mu x_* \tau_k] < (\varepsilon/2x^*). \qquad (2.26)$$

Consider the function $v\big(w(s), w(s + \tau_m - \tau_k)\big)$ for $s \in [t, t + \tau_k]$ defined by

$$v\big(w(s), w(s + \tau_m - \tau_k)\big) = \sum_{l=1}^{n} |\log w_l(s) - \log w_l(s + \tau_m - \tau_k)|. \qquad (2.27)$$

Using the fact $w$: $\boldsymbol{R} \to S$ and $w$ is a solution of (1.1), calculate the right derivative $D^+v \ of \ v$;

$$D^+v\big(w(s), w(s + \tau_m - \tau_k)\big) = \sum_{i=1}^{n} \{\text{sign}[w_l(s) - w_i(s + \tau_m - \tau_k)]\}.$$

$$\left\{\frac{dw_i(s)}{ds} \cdot \frac{1}{w_l(s)} - \frac{dw_l(s + \tau_m - \tau_k)}{ds} \frac{1}{w_l(s + \tau_m - \tau_k)}\right\}$$

$$\leq -\mu \sum_{l=1}^{n} |w_l(s) - w_l(s+\tau_m-\tau_k)| + \sum_{l=1}^{n} |b_i(s) - b_l(s+\tau_m-\tau_k)|$$

$$+x^* \sum_{l=1}^{n}\sum_{J=1}^{n} |a_{lJ}(s) - a_{lJ}(s+\tau_m-\tau_k)| \quad (2.28)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+ \sum_{i=1}^{n} |b_l(s) - b_l(s+\tau_m-\tau_k)| \quad (2.29)$$

$$+x^* \sum_{l=1}^{n}\sum_{j=1}^{n} |a_{lJ}(s) - a_{ij}(s+\tau_m-\tau_k)|.$$

If we set $s-\tau_k = \sigma$ then we have from (2.29),

$$D^+ v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+ \sum_{l=1}^{n} |b_l(\sigma+\tau_k) - b_l(s+\tau_m)|$$

$$+x^* \sum_{l=1}^{n}\sum_{j=1}^{n} |a_{ij}(\sigma+\tau_k) - a_{ij}(\sigma+\tau_m)| \quad (2.30)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+[(\varepsilon\mu x_*)/(4x^*)] + [(\varepsilon\mu x_*)/(4x^*)]$$

(2.31)

Integrating (2.31) over $(t, t+T_k)$

$$v(w(t+T_k)w(t+T_m))$$

Another MODEL where the model is directly applicable with the scheduled ones are the case of insurance and corporate finance. Insurance: Mathematics and Economics Volume 42, Issue 1, February 2008, Retrieval of Black–Scholes and generalized Erlang models by perturbed observations at a fixed time Dedicated to Professor André Dubey on the occasion of his 60th birthday Daniel Neuenschwander

S-stable laws on the real line (more generally on Hilbert spaces), associated with some non-linear transformations (so-called "shrinking operations"), were introduced in [Jurek, Z.J., 1977. Limit distributions for truncated random variables. In: Proc. 2nd Vilnius Conference on Probability and Statistics, June 28–July 3, 1977. In: Abstracts of Communications, vol. 3, pp. 95–96; Jurek, Z.J., 1979. Properties of s-stable distribution functions. Bull. Acad. Polon. Sci. Sér. Math. XXVII (1), 135–141; Jurek, Z.J., 1981. Limit distributions for sums of shrunken random variables. Dissertationes Math. vol. CLXXXV]. In [Jurek, Z.J., Neuenschwander, D., 1999. S-stable laws in insurance and finance and generalization to nilpotent Lie groups. J. Theoret. Probab. 12 (4), 1089–1107], the authors interpreted s-stable motions on the real line as limits of total amount of claims processes (up to a deterministic premium) of a portfolio of excess-of-loss reinsurance contracts and showed that they led to Erlang's model or to Brownian motion. In [Neuenschwander, D., 2000b. On option pricing in models driven by iterated integrals of Brownian motion. In: Mitt. SAV 2000, Heft 1, pp. 35–39], we considered stochastic integrals whose integrand and integrator are both independent Brownian motions, thus modelling a stochastic volatility; as a result we got an analogue of the Black–Scholes formula in this model, confirming a result of Hull and White [Hull, J., White, A., 1987. The pricing of options on assets with stochastic volatility. J. Finance XLII (2), 281–300]. In the present paper, we will look at a common generalization of these processes, namely s-stable motions on the real line perturbed by a stochastic integral whose integrand and integrator are both (not necessarily independent) s-stable motions. The main result will be that if we can observe the distribution of such so-perturbed s-stable motions (together with the values of the perturbing processes) at time t=1, then we can identify the whole model (including the perturbation) among all models with Lévy processes perturbed by an iterated stochastic integral of two Lévy processes (in the Gaussian case) resp. among all models with a compound Poisson process with drift perturbed by an iterated stochastic integral of two compound Poisson processes (in the completely non-Gaussian case if the perturbing processes have no drift) without knowing anything about the history or about its distribution during $0 \le t < 1$. This applies, e.g., to a situation

where several assets obey the same model and one can estimate the distribution at time one by looking at the values of all these assets at time t=1. Interestingly enough, it will be convenient to treat the whole matter in the algebraic framework of the so-called Heisenberg group. This is a concept coming in fact from quantum mechanics and is in a certain sense the simplest non-commutative Lie group.

$\leq e^{-\mu x^* T_k} v\big(w(t), w(t+T_m-T_k)\big) + (\varepsilon/2x^*)$

$\leq e^{-\mu x^* T_k} \sum_{i=1}^{n} |\log\ w_i(t) - \log\ w_i(t+T_m-T_k)| + (1/x_*) + (\varepsilon/2x^*)$

$\leq n(2x^*/x_*)\exp[-\mu x_* T_k] + (\varepsilon/2x^*)$

$<(\varepsilon/x^*)$.

Since

$$v\big(w(t+T_k)w(t+T_m)\big) \leq \sum_{i=1}^{n} |\ w_1(t+T_k) - w_1(t+T_m))/x^*)$$

We have from(2.33)that

$$\sum_{i=1}^{n} |\ w_1(t+T_k) - w_1(t+T_m)| < \varepsilon\ for\ m \geq k \geq n_0(\varepsilon)$$

Which show that the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\}$converges uniformly on **R**. By a similar argument the above conclusion is also valid when $T_k \to -\infty\ as\ k \to \infty$. Thus the uniform convergence of the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\}\ as\ k \to \infty\ on\ \mathbf{R}$ follows and this completes the proof of the theorem.

**COROLLARY.** Suppose $b_i a_{ij}$ (i, j = 1.2, … , n)in (1.1)are periodic with a common period say $T > 0$. *Then if* (1.2), (1.3)*and* (2.3)*hold the system* (1.1)*will have a periodic solution* $w(t) = \{w,(t), \ldots\ldots, w_n(t)\}$*of period* $T$ *such that*

$$w_1(t+T) = w_1(t), \qquad t \in \boldsymbol{R}: i = 1,2, \ldots, n.$$

**PROOF.** Proof follows from that of Theorem 2.1 if we Choose $T_k = KT(K = 0,1,2, \ldots)$*and note that*

$$x_1(t+KT) - w_1(t) \to 0\ as\ k \to \infty; i = 1,2, \ldots, n; t \in \boldsymbol{R},$$

$$x_i(t+(K+1)t) - w_1(t) \to 0 \ \ as \ \boldsymbol{K} \to \infty; i = 1,2,\dots,n; t \in \boldsymbol{R},$$

$$x_i(t+T+KT) - w_1(t+T) \to 0 \ \ as \ k \to \infty; i = 1,2,\dots,n; t \in \boldsymbol{R},$$

Implying that

$$w_1(t+T) = w_i(t); \quad i = 1,2,\dots,n; t \in \boldsymbol{R},$$

## 3. Global asymptotic stability

By definition we say that an almost-periodic solution say $w(t) = \{w_1(t),\dots,w_n(t)\} of (1.1) is\ globally\ asymptotically$ stable (or attractive) if and only if every other solution x(t)={ $x_1(t),\dots,x_n(t)$}with $x_1(o) > 0, i = 1.2,\dots,n$ is defined for all t□0 and satisfies

$$\lim_{t\to\infty} |w_1(t) - x_1(t)| = 0; \quad i = 1,2,\dots,n.$$

A consequence of such a global asymptotic stability is that there cannot be another strictly positive almost-periodic solution of (1.1)

**THEOREM 3.1** Suppose the almost- Periodic coefficients $b_1$, $a_{ij}$ (i, j = 1,2, … , n)

Satisfy (1.2)and (1.3)and the following

$$\inf_{t\in R} a_{ij}(t) > \left\{ \sum_{\substack{j=1 \\ j\neq 1}}^{n} \sup a_{ij}(t) \right\} + \mu; \qquad i = 1,2,\dots,n$$

For some positive number $\mu\ Then$ (1.1)has an almost-periodic solution (with strictly positive components) which is globally asymptotically stable.

We will omit the proof since the proof is identical with that of Theorem 3.1 of [3]

Under condition (3.2) the almost –periodic solution $w(t) = \{w_1(t), \ldots . w_n(t)\}$ of (1.1)enjoys a somewhat strong from of stability even if the coefficients of (1.1)are subjected to a class of perturbations We note (see Yoshizawa [7]page17)that if a function b:R→ $R$ is almost periodic then H(b),the "hull of b" is defined as follows:

H(b)= g|g:**R**→ $\boldsymbol{R}$ $and\ g(t)\lim_{k\to\infty} b(t+T_k) for$ some sequence

$T_k$ of real numbers the limit being uniform in $t \in R$.

By Theorem 2.2 of Yoshizawa [7, page 10], members of the hull of almost periodic functions are themselves almost periodic.

A system of the from

$$\frac{dy_1(t)}{dt} = y_1(t)\left\{B_i(t) - \sum_{J-1}^{n} A_{ij}\ (t)y_1(t)\right\},$$

*i=1,2,...,n; t>*$t_0 \in R$,

Where

$$B_i \in H(b_i); (A_{ij}), \qquad \text{i, j} = 1,2, \ldots, \text{n},$$

is known as a system from the hull of (1.1) As a consequence of the definition of the hull in (3.3)and (3.5)we shall have

$$\inf_{t\in R} B_i(t) = \text{b}_i^{\,l}; \quad \sup_{t\in R} B_i(t) = \text{b}_i^{\,l};$$

$$\inf_{t\in R} A_{ij}(t) = \text{b}_i^{\,l}; \quad \sup_{t\in R} A_{ij}(t) = \text{a}_{ij}^{\,u};$$

and hence the coefficients of(3.4)- (3.5)satisfy the same conditions as (1.2),(1.3)and (3.2)It will then follow from theorems(2.1)and (3.1)that solutions of (3.4)-(3.5)with $y_1(t_0) > 0, i = 1,2, \ldots, n\ will\ have$ the following property:

$$y_1(t) > 0 \ \ for\ t \geq t_0; i = 1,2,\dots,n,$$

$$\sum_{i=1}^{n} y_1(t) remains\ bounded\ for\ t \geq t_0$$

**DEFINITION** A Solution $x(t) = \{x_1(t), \dots, x_n(t)\}, t \in t_0 \in \boldsymbol{R}$,of (1.1)is said to be stable under perturbations from the hull of (1.1)if the following holds :let $y(t) = \{(y_1(t), \dots, y_n(t)\}$ be any solution of (3.4)-(3.5)such that $y_1(t_0) > 0$, i=1,2,…,n,and for some positive constants $\beta_1, \beta_2 \beta_1, \leq y_1(t) \leq \beta_2, t \geq t_0, \ \ i = 1,2,\dots,n$ Then given any $\varepsilon > 0$,there exists a constant $\delta > 0$ such that

$$\sum_{j=1}^{n} |x_1(t_0) - y_1(t_0)| < \delta$$

$$\sum_{i=1}^{n} \sup_{t \in R} |b_1(1) - B_1(t)| < \delta \quad = \sum_{i=1}^{n} |x_1(t) - y_i(t)| < \varepsilon$$

$$\sum_{i=1}^{n} \sum_{j=1}^{n} | \quad \sup_{t \in R} |a_{ij}\ (t) - A_{ij}(t)| < \delta$$

For $t \geq t_0$.

We can now formulate the following result which shows that all solution of (1.1)are stable under perturbations from the hull of (1.1)

**THEOREM 3.2**. Assume that the conditions of Theorem (3.1) hold. Then every strictly positive solution of (1.1) ( including its unique almost-periodic solution) is stable under perturbations from the hull of (1.1).

PROOF. Let $x(t) = \{x_1(t), \dots, x_n(t)\}$ and $y(t) = \{y_1(t), \dots, y_n(t)\}$ be any two strictly positive solutions of (1.1) and (3.4) – (3.5) respectively for $t \geq t_0 \in R$. Let $\beta_1, \beta_2$ be constants such that

$\beta_1 \leq x_i(t) \leq \beta_2$ ,
$\beta_1 \leq y_i(t) \leq \beta_2$ , $i = 1,2,\dots,n; t \geq t_0$ (3.7)

Consider now a function $V(t) = V(x(t)y(t))$ defined by

$$V(x(t)y(t)) = \sum_{i=1}^{n} |\log x_1(t) - \log, y_i(t)|, \quad t \geq t_0 \in \text{R}. \tag{3.8}$$

It is easy to see that

$$\frac{1}{\beta_2} \sum_{i=1}^{n} |x_i(t) - y_i(t)| \leq V(x(t)y(t))$$

$$\leq \frac{1}{\beta_1} \sum_{i=1}^{n} |x_i(t) - y_i(t)|, \; t \geq t_0 \in \text{R}. \tag{3.9}$$

Calculating the right derivative $D^+V$ of $V$ and using (1.1) and (3.4) we derive (after simplification) that

$$D^+V\left(x(t), y(t)\right) \leq -\mu \sum_{i=1}^{n} |x_i(t) - y_i(t)| + \sum_{i=1}^{n} |b_i(t) - B_i(t)|$$

$$+\beta_2 \sum_{j=1}^{n} \sum_{i=1}^{n} |a_{ij}(t) - A_{ij}(t)|; \quad t \geq t_0,$$

$$\leq -\mu\, \beta_1\, V\left(x(t), y(t)\right) + \sum_{i=1}^{n} |b_i(t) - B_i(t)|$$

$$+ \beta_2 \sum_{i=1}^{n} \sum_{j=1}^{n} |a_{ij}(t) - A_{ij}(t)|; \quad t \geq t_0, \tag{3.10}$$

An integration of (3.9) over $[t_0, t]$ with an application of (3.8) leads to

$$\frac{1}{\beta_2}\sum_{i=1}^{n}|x_i(t)-y_i(t)| \leq V(x(t),y(t)) \leq V(x(t_0),y(t_0)) + \left[\sum_{i=1}^{n}\sup_{t\in R}|b_i(t)-B_i(t)| + \beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}|a_{ij}(t)-A_{ij}(t)|\right]\frac{1}{\mu\beta_1},$$

$$t \geq t_0,$$

And hence

$$\sum_{i=1}^{n}|x_1(t)-y_1(t)| \leq \left(\frac{\beta_2}{\beta_1}\right)\sum_{i=1}^{n}|x_1(t_0)-y_1(t_0)| + \frac{\beta_2}{\mu\beta_2}\left[\sum_{i=1}^{n}\sup_{t\in R}|b_i(t)-B_i(t)|\right.$$

The construction reason is because is criticized in a number of usage guides: The reason for the long delays was because the costs greatly exceeded the original estimates. One objection to this construction is based on its redundancy: the word because (literally, by cause) contains within it the meaning of reason; thus saying the reason is because is like saying "The cause is by cause," which would never be said. A second objection is based on the claim that because can introduce only adverbial clauses and that reason is requires completion by a noun clause. Critics would substitute that for because in the offending construction: The reason for the long delays in completing the project was that the costs.

$$\left.+\beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}|a_{ij}(t)-A_{ij}(t)|\right]. \qquad (3.11)$$

Now for any $\varepsilon > 0$ if one chooses a $\delta > 0$ such that

$$\delta < \min\left\{\frac{\varepsilon}{3}\frac{\beta_1}{\beta_2}, \frac{\varepsilon}{3}\frac{\beta_1}{\beta^2}\mu\right\} \qquad (3.12)$$

Then (3..6) will be satisfied and this completes the proof

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# . Jia Li's Prescientiial Dynamics And Contential Staticity Of Accentuation Attrition Models Applied To Cosmological And Astrophysical And Quantum Mechanical Engineering Problems

**Abstract**: **Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics ET AL ). We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems. Predator prey model literally form bastion, pillar, post, stylobate, sentinel towards the end of dismasting many of the asymptotic and solutional behaviour in finance, economic, and as we shall see in cosmology, quantum computation and other fields.**

## MODEL FORMULATION

Let $p_1(t, a_1)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_1}{dt} = -b_1 Q_1 + g_1(g_1)Q_1$$

Where $b$ is the natural death rate in the absence of credit types and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_1(0) = 0, \quad g_1(P_1) > 0, \ P_1 \ 0, \qquad \lim_{P_1 \to \infty} g_1(P_1) = (K_1)_1 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_1 + (\rho_a)_1 + D_1(a_1, P_1, Q_1)\rho_1 + 0, \qquad 0 < a_1 \ \infty,$$

$$\rho_1(t, 0) = \int_0^{\infty} n_1(a_1, P_1, Q_1)\rho_1(t_1, a_1)da_1$$

Where $D_1$ is the mortality function and $n_1$ is the fertility function.

Put $_1$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_1}{dt}\right]_1 = x_1\left[\frac{dQ_2}{dt}\right]_1 - x_2\left[\frac{dQ_1}{dt}\right]_1$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_1 - x_4\left[\frac{dQ_2}{dt}\right]_1$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_1 - x_6\left[\frac{dQ_3}{dt}\right]_1$$

Where $x_1, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_1$ and $g_1(P_1)Q_1$

Now when you add the term $(b_1Q_1)$ with $g_1(P_1)(Q_1)$ we are infect see is that $(b_1Q_1)$ is credited incrementally part of $g_1(P_1)(Q_1)$ or instant nearly, when subtracted from $\frac{dQ_1}{dt}$, it means that $g_1(P_1)(Q_1)$ is slowly eaten by $\frac{dQ_1}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_1 = \mu_1(a_1) + (f_1)_1(P_1)_1 + (f_2)_1(Q)_1 \quad (2)$$

$$n_1 = \beta_1(a_1)h_1(Q_1), \quad (3)$$

Where $\mu_1$the age- dependent debit rate whit is $\mu_1(a_1)$ 0 $\forall a > 0$; $(f_1)_1$ is the density- dependent death rate with $f_1(0) = 0$, $(f_1^1)_1(P_1)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_1)_1(P_1) = (k_1)_1$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_1(0) = 0, (f_2^1)_1(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_1(Q_1) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^1(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_1 + (\rho_a)_1 = -\mu_1(a_1) + (f_1)_1P_1) + \big(f_2)_1(Q_1)\big)(\rho_1)_1 \quad (4)$$

$$\rho_1(t_1, 0) = h_1(Q_1)\int_0^\infty \beta_1(a_1)\rho_1(t, a_1)da, \quad (5)$$

$$\frac{dQ_1}{dt}$$

$$= -b_1Q_1 + g_1(P_1)Q_1 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_1$ and $Q_1$ of (4-6) satisfy

$$0 \leq \rho_1(t, a_1) \leq \rho_1^*(t, a_1),$$

$$0 \leq Q_1(t) \leq Q_1^*(t),$$

Where $\rho_1^*(t, a_1)$ and $Q_1^*(t)$, are solutions of

$$(\rho_t^*)_1 + (\rho_a^*)_1 = -(\mu_1(a_1) + (f_1)_1(P_1^*) + (f_2)_1 Q_1^*))\, \rho_1^*, \quad (7)$$

$$\rho_1^*(t, 0) = \int_0^\infty \beta_1(a_1)\, \rho_1^*(t, a_1) da_1 \quad (8)$$

$$\frac{dQ_1^*}{dt} = -b_1 Q_1^* + g_1(P_1^*) Q_1^*, \quad (9)$$

With $P_1^*(t) = \int_0^\infty \rho_1^*(t, a_1) da_1$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_1^*(t)}{dt} = P_1^*(t)\big(r_1 - (f_1)(P_1^*) - (f_2)_1(Q^*)\big), \quad (10)$$

$$\frac{dQ_1^*(t)}{dt} = Q_1^*(t)(-b_1 + g_1(P_1^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_1(a) e^{-\int_0^a \mu_1(\tau) dt} e^{-ra}\, da_1 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_1^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(t)dt}\, da_1 > 1, \qquad (H1)$$

Which guarantees $r_1 > 0$.

*Remark.* If $\beta_1(a) = [\beta_0]_1 a_1 e^{-\alpha_1 a_1}, \mu_1(a_1) \equiv \mu_1$ then the solution of (12) is

$$r_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_1]_1(P_1) = 0, [f_2](Q_1) = \lambda_1 Q_1\ h_1(Q_1) = 1, \qquad and\ \ g_1(P_1) = c_1 P_1$$

$$[\rho_1]_1 + [\rho_a]_1 = -(\mu_1 + \lambda_1 Q_1)\rho_1$$

$$\rho_1(t,0) = \int_0^{\infty} [\beta_0]_1 a_1 e^{-\alpha_1 a_1} \rho_1(t, a_1)\, da_1$$

$$\frac{dQ_1}{dt} = -b_1 Q_1 + c_1 P_1 Q_1$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_1}{dt} = \omega_1 P_1 - \lambda_1 P_1 Q_1$$

$$\frac{dQ_1}{dt} = b_1 Q_1 - c_1 P_1 Q_1$$

Where $\omega_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$ THE EQUILIBRIUM $(\rho_0)_1(a_1), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_1(a_1), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_1 + (\rho_a)_1 = -\big(\mu_1(a_1) + (f_1)_1(P_1)\big)\rho,$$

$$\rho_1(t,0) = \int_0^\infty \beta_1(a_1)\rho_1(t,a_1)da_1 \qquad (13)$$

The solution of (13) has the form of

$$\rho_1(t,a_1) = Ce_1^{r(t-a_1)} - \int_0^a (\tau)d\tau - \int_0^1 (f_1)_1\big(P_1(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_1(t) = C_1 e^{r_1 t - \int_0^t (f_1)_1(P_1(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_1^{-1}(r)$ if $k_1 > r$.

This is equivalent to

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{(k_1)_1 a_1} da_1 < 1.$$

In fact, if we define

$$R_1(s_1) = \int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau)d\tau} e^{-s_1 a_1}\, da_1$$

$R(s)$ is a decreasing function. $R_1(r_1) = 1$ And $R_1(k_1)_1 < 1$ imply $r_1 < (k_1)_1$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_1(a_1), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{k_1 a_1} da_1 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_1 \equiv \int_0^\infty (\rho_0)_{11}(a_1)da_1$ *satisfies*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{-(f_1)_1(P_0)a_1} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_1(a_1), 0)$

Let $u_1(t, a_1) = \rho(t, a_1) - \rho_0(a_1)$ Then

$$N_1(t) \equiv \int_0^\infty u_1(t, a_1)da_1 = \int_0^\infty \rho_1(t, a_1)da_1 - \int_0^\infty (\rho_0)_1(a_1)da = P_1(t) - P_0)_1$$

And

$$(u_1)_1 + (u_a)_1 = (\rho_t)_1 + (\rho_a)_1 - (\rho_0^1)_1(a_1)$$

$$= -\mu_1(a_1) + (f_1)_1(P_1) + (f_2)_1(Q_1))(u_1 + (\rho_0)_1) + \mu_1(a_1)(\rho_0)_1 + (f_1)_1(P_0)_1(\rho_0)_1$$

$$= -\mu_1(a_1)u_1 - (f_1)_1(P_1) + (f_2)_1(Q_1))u_1 - (\rho_0)_1(f_1)_1(P_1) - (f_1)_1(P_0)_1$$

$$+ (f_2)_1(Q_1)), \qquad (15)$$

Where $(\rho_0^1)_1(a_1)$ denotes $d(\rho_0)_1(a_1)/da_1$

Expansion $(f_1)_1(P_1)$ and $(f_2)_1(Q_1)$ at $(P_0)_1$ and 0 yield respectively

$$(f_1)_1(P_1) = (f_1)_1((P_0)_1 + N_1)$$

$$= (f_1)_1(P_0)_1 + (f_1^1)_1(P_0)_1N_1 + o_1(N_1), \qquad (16)$$

$$(f_2)_1(Q_1)$$

$$= (f_2^1)_1(0)Q + o_1(Q_1) \quad (17)$$

Write $f_1(P_0) = f_{10}, f_1'(P_0) = f_{10}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_1)_1 + (u_a)_1 \approx -(\mu_1(a_1) + (f_{10})_1)u - (f_{10}^1)_1(\rho_0)_1(N_1) - (f_{20}^1)_1(\rho_0)_1(Q_1)$$

$$u_1(t,0) = \rho_1(t,0)(\rho_0)_1(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t,a) + (\rho_0)_1(a_1)da_1 - (\rho_0)_1(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t,a_1)da_1 + (h_1(Q_1) - 1)(\rho_0)_1(0) \quad (18)$$

Expanding

$$h_1(Q_1) = 1 + (h^1)_1(0)Q_1 + 0(Q_1) \equiv 1 + h_0'Q_1 + 0Q_1$$

And substituting it into (18), we arrive at

$$u_1(t,0) \approx \int_0^\infty \beta_1(a_1)u_1(t,a_1)da_1 + (h_{a_1}')(\rho_0)_1Q_1$$

The linearization is the following:

$$[u_1] + [u_1] = -(\mu_1(a_1) + [t_{10}]_1u_1 - [\alpha_1]_1(a_1)N_1 - [\alpha_2]_1(a_1)Q_1$$

$$u_1(t,0) = \int_0^\infty \beta_1(a_1)u_1(t,a_1)da_1 + [\xi_1]Q_1 \quad (19)$$

$$\frac{dQ_1}{dt} = -b_1Q_1$$

where

$$[\alpha_1]_1(a_1) = [t_1^1 a]_1[p, o]_1(a_1)d2(a_1) = [f_{20}^1]_1(\rho\, o\,)_1(a_1)\xi = [\mathrm{h\ o}^1]_1[\rho \mathrm{o}]_1(\mathrm{o}) \text{ and } \mathrm{b}_1 = \mathrm{b}_1 - g_1(P_1)$$

If $g_1(P_0) > b$, clearly, $(\rho_0)_1(a_1), 0)$ is unstable and the predator population is persistent. Assume $g_1(P_0)_1 < b$.

Then, $Q_1(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_1 + u_a = -(\mu_1(a_1) + (f_{10})_1)u_1 - (\alpha_1)_1(a_1)N_1$$

$$= -(\mu_1(a_1) + f_{10})u_1 - \alpha_1(a_1)\int_0^\infty u_1(t, a_1)da_1$$

$$u_1(t, 0) = \int_0^\infty \beta_1\,(a_1)u_1(t, a_1)da_1$$

Using the method developed by Cushing [2, 3], we suppose

$$u_1(t, a_1) = \tilde{\mathrm{u}}_1(a_1)^{e^{\gamma(t-a_1)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_1}\, da_1$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_1)$ gives,

$$\tilde{\mathrm{u}}_1(a_1) = e^{-\int_0^a \mu_1(\tau_1)d\tau - (f_{10})_1 a}\left(\tilde{\mathrm{u}}_1(0) - w\int_0^a e^{\int_0^\tau \mu_1(s_1)ds + (f_{10})}\tau_1\,(\alpha_1)_1(\tau_1)e^{\gamma_1\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_1(a_1)e^{-\gamma_1 a_1 - (f_{10})_1 a_1 - \int_0^a \mu_1(\tau_1)d\tau}\right)\tilde{\mathrm{u}}_1(0)$$

$$+w\int_0^{\infty}\beta_1(a_1)e^{-\gamma_1 a_1-(f_{10})_1 a_1-\int_0^a \mu_1(\tau_1)d\tau}\int_0^a(\alpha_1)_1(\tau_1)e^{-\gamma_1\tau_1}e^{(f_{10})_1\tau_1+\int_0^{\tau}\mu_1(s_1)ds}d\tau_1\,da_1$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty}e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}da\right)u(0)$$

$$+\left(-1-\int_0^{\infty}e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}\int_0^a e^{f_{10}\tau+\int_0^{\tau}\mu(s)ds+y\tau}\alpha_1(\tau)d\tau\,da\right)w=0.\ (23)$$

Define $Y(\gamma,a)=e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta\,(a)Y(\gamma,a)da\right)u(0)$$

$$+w\int_0^{\infty}\beta(a)Y\,(\gamma,a)\int_0^a\alpha_1(\tau)Y^1(\gamma,\tau)d\tau\,da=0,(24)$$

$$\left(\int_0^{\infty}Y(\gamma,a)da\right)u(0)$$

$$+\left(-1-\int_0^{\infty}Y(\gamma,a)\int_0^a\alpha_1(\tau)Y^{-1}(\gamma,\tau)d\tau\,da\right)w=0.$$

Equation (20) has a solution of the form of $u(t,a)=u(a)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $u(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I(y,a)=\int_0^a\rho_0(\tau)Y^{-1}\,(\gamma,\tau)d\tau$. We obtain the characteristic equation

$$f_{10}^{l}\left(\int_0^\infty Y(\gamma,a)da\right)\left(\int_0^\infty \beta(a)Y(\gamma,a)I(\gamma,a)da\right)$$

$$+\left(1-\int_0^\infty \beta(a)Y(\gamma,a)da\right)\left(1+f_{10}^{l}\int_0^\infty Y(\gamma,a)I(\gamma,a)da\right)=0. \quad (25)$$

Since

$$\rho_0(a)=\rho_0(0)e^{-\int_0^a \mu(\tau)d\tau - f_{10}a}$$

$$I(\gamma,a)=\int_0^a \rho_0(\tau)e^{\gamma\tau+f_{10}\tau+\int_0^\tau \mu(s)ds}d\tau$$

$$=\frac{\rho_0(0)}{\gamma}(e^{\gamma a}-1)$$

Let

$$E(a)=e^{-f_{10}a-\int_0^a \mu(\tau)d\tau}$$

Then $[Y_1]_1[\gamma_1]_1,a_1)=\ E_1(a_1)e^{-\gamma_1 a_1}$. by a simple calculation,

$$\int_0^\infty \beta_1(a_1)Y_1(\gamma_1,a_1)I(\gamma_1,a_1)da=\frac{[\rho_0]_1(0)}{\gamma}\left(1-\int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}\ da\right);$$

$$\int_0^\infty Y_1(\gamma_1,a_1)I(\gamma_1,a_1)da=\frac{[\rho_0]_1(0)}{\gamma_1}\left(\int_0^\infty E_1(a_1)e^{-\gamma_1 a_1}\ da_1\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_1(0)}{\gamma_1}[f_{10}^{1}]\int_0^\infty E_1(a_1)e^{-\gamma_1 a_1}da_1\ (\int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1)$$

$$+\left(1-\int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right).$$

$$\left(1+\frac{(\rho_0)_1(0)}{\gamma_1}(f_{10}^{'})_1\left(\int_0^\infty E_1(a_1)da_1-\int_0^\infty E_1(a_1)e^{-\gamma_1 a_1}da_1\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right)(1+\frac{(f_{10}^{'})_1}{\gamma_1}(\rho_0)_1(0)\int_0^\infty E_1(a_1)da_1)=0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_1(a_1)Y_1(\gamma_1 a_1)da_1 = 1-\int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1 \neq 0,$$

Since $\int_0^\infty \beta_1(a_1)Y_1(\gamma_1 a_1)\int_0^a \alpha_1(\tau_1)Y^{-1}(\gamma_1,\tau_1)d\tau_1 da_1 > 0$. Therefore, the characteristic solution is

$$\gamma_1 = -(f_{10}^{'})_1(\rho_0)_1\int_0^\infty E_1(a_1)da_1.$$

Since $f_1'(P_0)_1 > 0$, we obtain the following result

THEOREM 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_1), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_1), 0)$ which is related to the extinction of debit summation as has been referred to in the foregoing and to The possible persistence of the prey

population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_1 + (\mu_1(a_1) + [f_1]_1(P_1^*) + [f_2]_1(Q_1^*))P_1^* = 0, \quad (26)$$

$$p_1^*(0) = h_1(Q_1^*)\int_0^\infty \beta_1(a_1)p_1^*(a_J)p_1^*(a_1)da_1 \quad (27)$$

$$g_1(P_1^*) - b_1 = 0, \quad (28)$$

Where $P_1^* = \int_0^\infty p_1^*(a_1)\, da_1$

Equation (28) yields $P^* = g^{-1}(b)$,

From (26),

$$p_1^*(a) = p_1^*(0)e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}$$

Substituting into (27), we need to solve

$$\left(1 - h_1(Q_1^*)\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}da_1\right)p_1^*(0) = 0$$

For $Q_1^*$

Fix $p_1^*(0)$ and $P_1^*$ Let

$$H_1(z_1) = 1 - h_1(z_1)\int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1}da_{,1}$$

And assume that

$$[k_3]_1 \int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(g_1^{-1}(b_1))a_1 - [k_2]_1 a_1}\, da_1 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty} [H_1(z_1)] > 0.$$

Clearly,

$$\frac{dH_1}{dz_1} = h_1(z)([f_2^{'}]_1(z_1)) \int_0^{\infty} a_1\beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - f_1(P_1^*)a_1 - f_2(z_1)a_1} da_1$$

$-h^{'}(z_1) \int_0^{\infty} \beta_1(a_1)\, e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1} \quad da_1$

By the assumptions on $\left[f_2^1\right]_1 and\ [h^1]^1 dH_1/dz > 0$. Hence, if

$$H_1(0) = 1 - \int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1} da_1 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_1(0) < 0$ follows if $g_1^-(b_1) < [f_1^{-1}]_1(r_1)$, where $r_1$ is the solution of (12).

Indeed, since

$$R_1(s_1) = \int_0^{\infty} \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 e^{-y_1^a}} da_1$$

Is a decreasing function of $s_1$ when $P_1^* = [g^{-1}]_1(b_1) < [f_1^{-1}]_1(r_1)$, that is $[f_1]_1(P_1^*) < r$, $R\big(f_1(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_1(g^{-1}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1\left(g_1^{-1}(b)\right)a_1} > 1. \qquad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_1^*(a_1), Q_1^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_1^*(a_1), Q_1^*)$ is more complicated. We again consider perturbations $m_1(t_1 a_1) = p_1(t_1, a_1) - p_1^*(a_1)$ and $W_1(t) = Q_1(t) - Q^*$.Then

$M_1(t) \equiv \int_0^\infty m_1(t, a_1)da_1 = P_1(t) - P_1^*,$

$$[m_1]_1 + [m_a]_1 = [p_1]_1 + [p_a]_1 - \left[p^*\right]_1(a_1) \approx -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1$$

$$-[f_{1*}^1]_1 p_1^*(a_1)M - [f_{2*}^1]_1 p_1^*(\mathrm{a})W_1,$$

$$m_1(t,0)p_1(t,0) - p_1^*(0) \approx h_* \int_0^\infty \beta_1(a_1)m_1(t,a_1)\,da + \frac{[h_*^1]_1}{[h_*]_1}\, p_1^*(0)W_1$$

$$\frac{dW_1}{dt} = \frac{dQ_1}{dt} \approx [g_*^1]_1 Q_1^* M_{,1}$$

Where $[f_{1*}]_1$ denoting $[f_1(P^*)]_1$, $[f_{1*}]_1$ denoting $[f_2(P^*)]_1$, $\left[f_{1*}^{'}\right]_1$ denoting $\left[f^{'}{}_1(P^*)\right]_1 \left[f_{2*}^{'}\right]_1$ denoting

$[f_2]_1\,(Q_1^*,\quad h_{1*}^{'}$ denoting $h_*^{'}$ denoting $h_1^{'}\,(Q_1^*)$, and $\left[g_*^{'}\right]_1$

Denoting $\left[g^{'}\right]_1(P_1^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_1]_1 + [m_a]_1 = -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1 - [f'_{1*}]_1[p_1^*](a_1)M_1 - [f'_{2*}]_1 p_1^*(a_1)W_1$$

$$m_1(t,0) = [h_*]_1 \int_0^\infty \beta_1(a_1)m_1(t,a_1)da_1 + \frac{[h_*^1]_1}{[h^*]_1}[p^*]_1(0)\,W_{,1} \qquad (30)$$

$$\frac{dW_1}{dt} = [g_*^1]_1[Q_1^*]M_1$$

Integrating $[m_1]_1(t,a_1)$ with respect to a from 0 to $\infty$, we have

$$\frac{dM_1(t)}{dt} = \int_0^\infty [mt]_1(t,a_1)da_1$$

$$= -\int_0^\infty [m_a]_1\,(t,a_1)da_1 - \int_0^\infty \mu(a_1)m(t,a_1)da_1$$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f'_{1*}]_1 \int_0^\infty p_1^*(a_1)da_1\right)M_1 - [f_{2*}^1]_1 \int_0^\infty p_1^*(a_1)da_1 W_1$$

$$= m_1(t,0) - \int_0^\infty \mu_1(a_1)\,m_1(t,a_1)da_1$$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f_{1*}']_1 P_1^*\right)M_1 - [f_{2*}']_1 P_1^* W_{,1}$$

From (30), we arrive at the equations

$$\frac{dW_1(t)}{dt} = \int_0^\infty ([h_*]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1$$

$$-\left([f_{1*}] + [f_{2*}]_1 + [f_{1*}']_1 P_1^*\right)M_1 + \left(\frac{[h_*']_1}{[h_*]^1}\,P_1^*(0) - f_{2*}' P_1^*\right)W_1,$$

(31)

$$\frac{dW_1(t)}{dt}\left[g^{'}{}_{*}\right]_1 Q_1^{*} M_{.1}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{*}]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1 = ([h_{*}]_1\beta_1(\theta_1) - \mu_1(\theta_1))M_{,1}$$

Where $0< \theta_1 = \theta_1(t) < \infty$.

And if $a_1^2\,[a_1^2]_1<4\,[a_2]_1<4\,a_2 e$,

Re $\lambda_1\left[\lambda_1\right]_1 = \text{Re}\lambda_2\left[\lambda_2\right]_1 = -\frac{1}{2}a_1[a2\ ]_1 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_{*}\,[h_{*}]_1\beta_1(a_1) - \mu_1(a_1)\right|$. Then$|n_1(t)|$ ⊠Δ∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium($\rho^{*}(a), \boldsymbol{Q}^{*}$) is asymptotically stable if either $a_1^2 <4\,a_2 e$ and $[a_1] > 2\Delta$,or $[a_1^2]$ ⊠4$[a_2]$e and $(a_1 - \sqrt{(a_1^2)_1 - 4(a_2)_1 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+1}\right)$ and $\beta_1(a_1) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1) + d\tau_1 da} = \frac{1}{(m+\mu_0+1)^2} + \frac{2}{(m+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_1 + [\mu_0]$ 0.5. Assume$[f_1](x_1) = K_1 x_1$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_1(x_1) = [0x_1]/(1+x_1)$; the functional responses $[[f_2](x) = [0]\, x_1/(1+x_1)$and $h_1(x_1)$=$(1+\omega_1 x_1)/(1+x_1)$(See Holling [7] and Rashevsky [13])

Equation (14),now has the form of

$$\frac{m_1 + [\mu_0] + [K][P_0])}{[m_1]_1 + (\mu_0)_1 + [K][P_0])_1^3\ [K]_1 [P_0]_1^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1+P_0) < b_1$i e $[P_0] < b/(\delta_1 - b_1)$then the equilibrium$(P_0)$ $[P_0]_1(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_1$

$[f_{1*}] = K_1\ P^*\ [f_{2*}] = \frac{\sigma Q_1^*}{1+Q_1^*} \quad [h_*] = \frac{1+\omega_1 Q_1^*}{1+Q_1^*}$

$[f_{1*}] = K_1\ [f_{2*}^1] = \frac{\sigma_1}{(1+Q_1^*)^2}[g_*^1] = \frac{\sigma_1}{(1+P_1^*)^2} \quad [h_*^1] = \frac{\omega_1 - 1}{(1+Q_1^*)^2}$

## AN ANALOGY BETWEEN PREDATOR PREY AND ACCENTUATION ATTRION MODELS

**Best example how the model could be applied is the Schrödinger's equation itself wherein RHS gobbles up LHS, and the terms on RHS do the same with each other. It is like adding milk to water. Whether it happens instantaneously or time lag is there is a different question altogether dealt at the opportune points**

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| **Best example that can be given is the formulation of the quantum wave equation albeit in a nascent form for the recently found gravitational waves. This has been studied in length in various papers of gravitational waves, wherein the attribution of how the model is used is described. Note the calculation holds for only one module. This is in consideration to the fact that the relation is beteen the both debit systems and credit systems which as we have enucleated in the foregoing bears ample testimony for the application of the model** |
| **Gravitational quantum wave equation:** |
| **Equation Time-dependent equation (Wikipedia)** |
| **(1) In quantum mechanics, the Schrödinger equation is (=) a partial differential equation that describes how the quantum state of a quantum system changes with (e&eb) time. It was formulated in late 1925, and published in 1926, by the Austrian physicist Erwin Schrödinger.[1]** |
| **(2) In classical mechanics Newton's second law, (F = ma), is used to make (eb) a mathematical prediction as to what path a given system will take following a set of known initial conditions.** |
| **(3) In quantum mechanics, the analogue of Newton's law is (=) Schrödinger's equation for (e) a quantum system (usually atoms, molecules, and subatomic particles whether free, bound, or localised).** |
| **(4) It is not a simple algebraic equation, but in general a linear partial differential equation, describing (eb) the time-evolution of the system's wave function (also called a "state function").[2]:1–2** |
| **(5) The concept of a wavefunction is (=) a fundamental postulate of quantum mechanics.** |
| **(6) Although Schrödinger's equation is often presented as a separate postulate, some authors [3]: Chapter 3 show that some properties resulting from Schrödinger's equation may be deduced just from (e) symmetry principles alone, for example the commutation relations.** |
| **(7) Generally, "derivations" of the Schrödinger equation demonstrate (eb) its mathematical plausibility for describing (eb) wave–particle duality, but to date there are no universally accepted derivations of (e) Schrödinger's equation from (e) appropriate axioms.** |
| **(8) In the Copenhagen interpretation of quantum mechanics, the wave function is (=) the most complete description that can be given of (e) a physical system.** |
| **(9) Solutions to Schrödinger's equation describe (eb) not only molecular, atomic, and subatomic systems, but also macroscopic systems, possibly even the whole universe.[4]:292ff** |

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| **(10) The Schrödinger equation is consistent with (e&eb) classical mechanics but not with (e&eb) special relativity.** |
| **(11) Making quantum mechanics consistent with (e&eb) special relativity is one of the great achievements of (e) Quantum Field Theory and requires (e) the introduction of the creation and annihilation operators so that (e) the number of particles can no longer be considered constant.** |
| **(12) The Schrödinger equation is not the only way to make predictions in (eb) quantum mechanics—other formulations can be used, such as Werner Heisenberg's matrix mechanics, and Richard Feynman's path integral formulation** |
| **(13) The form of the Schrödinger equation depends on (=) the physical situation (see below for special cases).** |
| **(14) The most general form is the time-dependent Schrödinger equation, which gives a description of a system evolving with time: [5]:143** |
| **(15) A wave function that satisfies the non-relativistic Schrödinger equation with V = 0. In other words, this corresponds to (e &eb) a particle traveling freely through (e&eb) empty space. Time-dependent** |
| **Schrödinger equation (general)** |
| **(16) i ℏ ∂ ∂ t Ψ ( r , t ) = H ^ Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t) where i is the imaginary unit, ħ is the Planck constant divided by 2π, the symbol ∂/∂t indicates a partial derivative with respect to time t, Ψ (the Greek letter psi) is the wave function of (e) the quantum system, and Ĥ is the Hamiltonian operator (which characterizes the total energy of any given wave function and takes different forms depending on the situation).** |
| **(17) Each of these three rows is a wave function which satisfies (eb) the time-dependent Schrödinger equation for a harmonic oscillator. Left: The real part (blue) and imaginary part (red) of the wave function. Right:** |
| **(18) The probability distribution of finding the particle with (e&eb) this wave function at a given position.** |
| **(19) The top two rows are examples of stationary states, which correspond to (e&eb) standing waves.** |
| **(20) The bottom row is an example of a state which is (=) not a stationary state.** |
| **(21) The right column illustrates (eb) why stationary states are (=) called "stationary".** |

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| **(22) The most famous example is the non-relativistic Schrödinger equation for (e) a single particle moving in an electric field (but not a magnetic field; see the Pauli equation): [6]** |
| **Time-dependent Schrödinger equation** |
| **(Single non-relativistic particle)** |
| **(23) i ℏ ∂ ∂ t Ψ ( r , t ) = [ − ℏ 2 2 μ ∇ 2 + V ( r , t ) ] Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)** |
| **(24) Here μ is (=) the particle's "reduced mass", V is its potential energy, ∇2 is the Laplacian (a differential operator), and Ψ is (=) the wave function (more precisely, in this context, it is called the (=) "position-space wave function").** |
| **(25) In plain language, it means "total energy equals (=) kinetic energy plus potential energy", but the terms take unfamiliar forms for (e) reasons explained below.** |
| **(26) Given the particular differential operators involved, this is (=) a linear partial differential equation.** |
| **(27) It is also a diffusion equation, but unlike the heat equation, this one is (=) also a wave equation given (e&eb) the imaginary unit present in the transient term.** |
| **(28) The term "Schrödinger equation" can refer to (e) both the general equation (first box above), or the specific nonrelativistic version (second box above and variations thereof).** |
| **(29) The general equation is indeed quite general, used throughout quantum mechanics, for everything from the Dirac equation to quantum field theory, by plugging in various complicated expressions for the Hamiltonian. The specific nonrelativistic version is a simplified approximation to reality, which is quite accurate in many situations, but very inaccurate in others (see relativistic quantum mechanics and relativistic quantum field theory)** |
| **(30) To apply the Schrödinger equation, the Hamiltonian operator is set up for the system, accounting for the kinetic and potential energy of the particles constituting the system, then inserted into the Schrödinger equation. The resulting partial differential equation is solved for the wave function, which contains information about the system.** |
| **Time-independent equation** |

**(31) The time-dependent Schrödinger equation described above predicts that wave functions can form standing waves, called stationary states (also called "orbitals", as in atomic orbitals or molecular orbitals). These states are important in their own right, and if the stationary states are classified and understood, then it becomes easier to solve the time-dependent Schrödinger equation for any state. Stationary states can also be described by a simpler form of the Schrödinger equation, the time-independent Schrödinger equation. (This is only used when the Hamiltonian itself is not dependent on time explicitly. However, even in this case the total wave function still has a time dependency.)**

**Time-independent Schrödinger equation (general)**

**(32) E Ψ = H ^ ⊠ Ψ {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi } {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi }**

**In words, the equation states:**

**(33) When the Hamiltonian operator acts on a certain wave function Ψ, and the result is proportional to the same wave function Ψ, then Ψ is a stationary state, and the proportionality constant, E, is the energy of the state Ψ.**

**(34) The time-independent Schrödinger equation is discussed further below. In linear algebra terminology, this equation is an eigenvalue equation.**

**(35) As before, the most famous manifestation is the non-relativistic Schrödinger equation for a single particle moving in an electric field (but not a magnetic field):**

**Time-independent Schrödinger equation (single non-relativistic particle)**

**(36) E Ψ ( r ) = [ − ℏ 2 2 μ ∇ 2 + V ( r ) ] Ψ ( r ) {\displaystyle E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )} E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )**

**Gravitational waves: (Wikipedia)**

**(37) In theory, the loss of energy through gravitational radiation could eventually drop the Earth into the Sun. However, the total energy of the Earth orbiting the Sun (kinetic energy + gravitational potential energy) is about 1.14×1036 joules of which only 200 Watts (joules per second) is lost through gravitational radiation, leading to a decay in the orbit by about 1×10−15 meters per day or roughly the diameter of a proton. At this rate, it would take the Earth approximately 1×1013 times more than the current age of the Universe to spiral onto the Sun. This estimate overlooks the decrease in r over time, but the majority of the time the bodies are far apart and only radiating slowly, so the difference is unimportant in this example.**

**More generally, the rate of orbital decay can be approximated by[39]**

**(38)** $\frac{dr}{dt} = -\frac{64}{5}\frac{G^3}{c^5}\frac{(m_1 m_2)(m_1 + m_2)}{r^3}$ , {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,} {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,}

**where r is the separation between the bodies, t time, G Newton's constant, c the speed of light, and m1 and m2 the masses of the bodies. This leads to an expected time to merger of [39]**

$t = \frac{5}{256}\frac{c^5}{G^3}\frac{r^4}{(m_1 m_2)(m_1 + m_2)}$. {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.} {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.}

**(39) For example, a pair of solar mass neutron stars in a circular orbit at a separation of 1.89×108 m (189,000 km) has an orbital period of 1,000 seconds, and an expected lifetime of 1.30×1013 seconds or about 414,000 years. Such a system could be observed by LISA if it were not too far away. A far greater number of white dwarf binaries exist with orbital periods in this range. White dwarf binaries have masses in the order of the Sun, and diameters in the order of the Earth. They cannot get much closer together than 10,000 km before they will merge and explode in a supernova which would also end the emission of gravitational waves. Until then, their gravitational radiation would be comparable to that of a neutron star binary.**

**(40) By substituting value of dr/dt in Schrödinger's wave equation we get the quantum wave equation for gravitational waves.**

In order to have (H3) and (H4) hold, $m, \mu_0$, K, b, δ, and $\omega$ need to satisfy

$$\frac{\omega_1}{(m + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 + b_1} + \sigma)^2} + \frac{2\omega_1}{([m_1]_1 + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 + b_1} + \sigma_1)^3} < 1$$

And

$$\frac{1}{(m_1 + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 - b_1})^2} + \frac{2\omega_1}{(m_1 + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 - b_1})^3} > 1$$

When the inequalities (33) and (34) hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$(m_1 + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 - b_1} + (m_1 + [\mu_0]_1 + \frac{K_1 b_1}{\delta_1 - b_1} + \sigma_1)\, Q_1^*)^3$

$= m_1 + [\mu_0]_1 + 3 + \frac{K_1 b_1}{\delta_1 - b_1}(1 + Q^*)(1 + \omega_1 Q_1^*$

Now

$$\left[a_1^2\right] - 4\left[a_2\right] e = \left(\frac{2K_1 b_1}{\delta_1 - b_1}\right)^2 + \frac{4K_1 b_1 \sigma_1 Q_1^*}{(\delta_1 - b_1)(1 + Q_1^*)} + \frac{\sigma_1^2 Q_1^{*2}}{(1 + Q_1^*)^2}$$

$$\frac{4\, b_1 \sigma_1 (\delta_1 - b) Q^*}{\delta_1 (1 + Q_1^*)^2} - \frac{4\, (\delta_1 - b_1)^2 (1 - \omega_1) Q_1^*}{\delta_1 (1 + \omega_1 Q_1^*)(1 + Q_1^*)} P^*(0)$$

And

$\Delta = \lim_{0 < a < \infty} |[h_*]_1\, a e^{mu}_1\ \ [\mu_0]_1 - \frac{a_1}{a_1 + 1}$

Hence if $[a_1^2 > 4(a_2)_1 e$ and

$$[a_1] = \frac{2K_1 b_1}{\delta_1 - b_1} + \frac{\sigma_1 Q_1^*}{1 + Q_1^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_1]_1 = [f_{1*}]_1 + [f_{2*}]_1 + [f_{1*}] P_1^*$$

$$[a_2]_1 = \left[f^{1}_{2*}\right] P^{*}_{1} - \frac{h^{1}_{*}}{h_{*}} P^{*}(0)$$

$$e=[g^{1}_{*}]_1 Q^{*}_{1}$$

$$\eta_1(t) = \left[h_{*}\right]_1 \beta_1 \left(\theta_1(t)\right) - \mu_1(\theta(t)),$$

The system(31)can be written as the following system:

$$[M(t)]\left[a_1\right]M_1 - [a_2]_1 W_1 - [n(t)]M_1$$

$$w_1(t) = eM_1$$

Write

$$A=\begin{pmatrix} -[a_1]_1 - [a_2]_1 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_1 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_1 \\ W_1 \end{pmatrix}$$

The solution of (32)is

$$X_1(t) = e^{A,t}[X]_1(0) + \int_0 e^{At}$$

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# Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

## (Module Three)

**Abstract: Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-435) 1990. We demonstrate its applicability cosmology/ quantum computation economics et al. As stated in the foregoing we shall apply model to various cosmological problem and show the relevance to the same. . To exemplify model can be applied to the following situational significations: Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an**

**outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 Masaru Shibata and Takashi Nakamura** perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. **Authors** use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society **Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995**

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(g_3)Q_3$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_3(0) = 0, \quad g_3(P_3) > 0, \quad P_3 \text{ 0}, \qquad \lim_{P_3 \to \infty} g_3(P_3) = (K_3)_3 \text{ ∞}. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_3 + (\rho_a)_3 + D_3(a_3, P_3, Q_3)\rho_3 + 0, \qquad 0 < a_3 \text{ ∞},$$

$$\rho_3(t,0) = \int_0^\infty n_3(a_3, P_3, Q_3)\rho_3(t_3, a_3)da_3$$

Where $D_3$ is the mortality function and $n_3$ is the fertility function.

Put $_3$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_3}{dt}\right]_3 = x_3\left[\frac{dQ_2}{dt}\right]_3 - x_2\left[\frac{dQ_3}{dt}\right]_3$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_3 - x_4\left[\frac{dQ_2}{dt}\right]_3$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_3 - x_6\left[\frac{dQ_3}{dt}\right]_3$$

Where $x_3, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_3$ and $g_3(P_3)Q_3$

Now when you add the term $(b_3Q_3)$ with $g_3(P_3)(Q_3)$ we are infect see is that $(b_3Q_3)$ is credited incrementally part of $g_3(P_3)(Q_3)$ or instant nearly, when subtracted from $\frac{dQ_3}{dt}$ , it means that $g_3(P_3)(Q_3)$ is slowly eaten by $\frac{dQ_3}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_3 = \mu_3(a_3) + (f_3)_3(P_3)_3 + (f_2)_3(Q)_3 \quad (2)$$

$$n_3 = \beta_3(a_3)h_3(Q_3), \quad (3)$$

Where $\mu_3$the age- dependent debit rate whit is $\mu_3(a_3)$ 0 $\forall a > 0$; $(f_3)_3$ is the density- dependent death rate with $f_3(0) = 0$, $(f_3^3)_3(P_3)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_3)_3(P_3) = (k_3)_3$ ; $f_2$ is the functional response of

creditors to prey with $[f_2]_3(0) = 0, (f_2^3)_3(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_3\ (Q_3) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 3, h^3(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_3 + (\rho_a)_3 = -\mu_3(a_3) + (f_3)_3 P_3) + \big(f_2)_3(Q_3)\big)(\rho_3)_3 \quad (4)$$

$$\rho_3(t_3, 0) = h_3(Q_3)\int_0^\infty \beta_3\,(a_3)\rho_3(t, a_3)da, \quad (5)$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(P_3)Q_3 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_3$ and $Q_3$ of (4-6) satisfy

$$0 \ \rho_3(t, a_3) \ \rho_3^*(t, a_3),$$

$$0 \ Q_3(t) \ Q_3^*(t),$$

Where $\rho_3^*(t, a_3)$ and $Q_3^*(t)$, are solutions of

$$(\rho_t^*)_3 + (\rho_a^*)_3 = -(\mu_3(a_3) + (f_3)_3(P_3^*) + (f_2)_3 Q_3^*))\,\rho_3^*, \quad (7)$$

$$\rho_3^*\,(t, 0) = \int_0^\infty \beta_3(a_3)\,\rho_3^*\,(t, a_3)da_3 \quad (8)$$

$$\frac{dQ_3^*}{dt} = -\,b_3\,Q_3^* + g_3(P_3^*)Q_3^*, \quad (9)$$

With $P_3^*(t) = \int_0^\infty \rho_3^*\ (t, a_3)da_3$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_3^*(t)}{dt} = P_3^*(t)\big(r_3 - (f_3)(P_3^*) - (f_2)_3(Q^*)\big), \quad (10)$$

$$\frac{dQ_3^*(t)}{dt} = Q_3^*(t)(-b_3 + g_3(P_3^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_3(a)e^{-\int_0^a \mu_3(\tau)dt}e^{-ra}\ da_3 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_3^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(t)dt}\ da_3 > 1, \quad (H1)$$

Which guarantees $r_3 > 0$.

*Remark.* If $\beta_3(a) = [\beta_0]_3 a_3 e^{-\alpha_3 a_3}, \mu_3(a_3) \equiv \mu_3$ then the solution of (12) is

$$r_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_3]_3(P_3) = 0, [f_2](Q_3) = \lambda_3 Q_3\ h_3(Q_3) = 1, \quad and\ \ g_3(P_3) = c_3 P_3$$

$$[\rho_3]_3 + [\rho_a]_3 = -(\mu_3 + \lambda_3 Q_3)\rho_3$$

$$\rho_3(t,0) = \int_0^{\infty} [\beta_0]_3 a_3 e^{-\alpha_3 a_3} \rho_3(t,a_3)\, da_3$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + c_3 P_3 Q_3$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_3}{dt} = \omega_3 P_3 - \lambda_3 P_3 Q_3$$

$$\frac{dQ_3}{dt} = b_3 Q_3 - c_3 P_3 Q_3$$

Where $\omega_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$

## THE EQUILIBRIUM $(\rho_0)_3(a_3), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_3(a_3), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_3 + (\rho_a)_3 = -\big(\mu_3(a_3) + (f_3)_3(P_3)\big)\rho,$$

$$\rho_3(t,0)$$

$$= \int_0^\infty \beta_3(a_3)\rho_3(t,a_3)da_3 \qquad (13)$$

The solution of (13) has the form of

$$\rho_3(t,a_3) = Ce_3^{r(t-a_3)} - \int_0^a (\tau)d\tau - \int_0^3 (f_3)_3\big(P_3(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_3(t) = C_3 e^{r_3 t - \int_0^t (f_3)_3(P_3(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_3^{-3}(r)$ if $k_3 > r$. This is equivalent to

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{(k_3)_3 a_3} da_3 < 1.$$

In fact, if we define

$$R_3(s_3) = \int_0^\infty \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau)d\tau} e^{-s_3 a_3}\, da_3$$

$R(s)$ is a decreasing function. $R_3(r_3) = 1$ And $R_3(k_3)_3 < 1$ imply $r_3 < (k_3)_3$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_3(a_3), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{k_3 a_3} da_3$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_3 \equiv \int_0^\infty (\rho_0)_{33}(a_3)da_3$ *satisfies*

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{-(f_3)_3(P_0)a_3} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_3(a_3), 0)$

Let $u_3(t, a_3) = \rho(t, a_3) - \rho_0(a_3)$ Then

$$N_3(t) \equiv \int_0^{\infty} u_3(t, a_3) da_3 = \int_0^{\infty} \rho_3(t, a_3) da_3 - \int_0^{\infty} (\rho_0)_3(a_3) da = P_3(t) - P_0)_3$$

And

$$(u_3)_3 + (u_a)_3 = (\rho_t)_3 + (\rho_a)_3 - (\rho_0^3)_3(a_3)$$

$$= -\mu_3(a_3) + (f_3)_3(P_3) + (f_2)_3(Q_3))(u_3 + (\rho_0)_3) + \mu_3(a_3)(\rho_0)_3 + (f_3)_3(P_0)_3(\rho_0)_3$$

$$= -\mu_3(a_3)u_3 - (f_3)_3(P_3) + (f_2)_3(Q_3))u_3 - (\rho_0)_3 (f_3)_3(P_3) - (f_3)_3(P_0)_3$$

$$+ (f_2)_3(Q_3)), \qquad (15)$$

Where $(\rho_0^3)_3(a_3)$ denotes $d(\rho_0)_3(a_3)/da_3$

Expansion $(f_3)_3(P_3)$ and $(f_2)_3(Q_3)$ at $(P_0)_3$ and 0 yield respectively

$$(f_3)_3(P_3) = (f_3)_3((P_0)_3 + N_3)$$

$$= (f_3)_3(P_0)_3 + (f_3^3)_3(P_0)_3 N_3 + o_3(N_3), \qquad (16)$$

$$(f_2)_3(Q_3)$$

$$= (f_2^3)_3(0)Q + o_3(Q_3) \qquad (17)$$

Write $f_3(P_0) = f_{30}, f_3'(P_0) = f_{30}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_3)_3 + (u_a)_3 \approx -(\mu_3(a_3) + (f_{30})_3)u - (f_{30}^3)_3(\rho_0)_3(N_3) - (f_{20}^3)_3(\rho_0)_3(Q_3)$$

$$u_3(t, 0) = \rho_3(t, 0)(\rho_0)_3(0)$$

$$= h_3(Q_3)\int_0^{\infty} \beta_3(a_3)(u_3(t,a) + (\rho_0)_3(a_3)da_3 - (\rho_0)_3\,(0)$$

$$= h_3(Q_3)\int_0^{\infty} \beta_3(a_3)(u_3(t,a_3)da_3 + (h_3(Q_3) - 3)(\rho_0)_3\,(0) \qquad (18)$$

Expanding

$$h_3(Q_3) = 1 + \ (h^3)_3(0)Q_3 + 0(Q_3) \equiv 1 + h_0^{'}Q_3 + 0Q_3$$

And substituting it into (38), we arrive at

$$u_3(t,0) \approx \int_0^{\infty} \beta_3\ (a_3)u_3(t,a_3)da_3 + (h_{a_3}^{'})(\rho_0)_3Q_3$$

The linearization is the following :

$$[u_3] + [u_3] = -(\mu_3(a_3) + [t_{30}]_3u_3 - [\alpha_3]_3(a_3)N_3 - [\alpha_2]_3(a_3)Q_3$$

$$u_3(t,0) = \int_0^{\infty} \beta_3(a_3)u_3(t,a_3)da_3 + [\xi_3]Q_3 \qquad (19)$$

$$\frac{dQ_3}{dt} = -b_3Q_3$$

where

$$[\alpha_3]_3(a_3) = [t_3^3a]_3[p,o]_3(a_3)d2(a_3) = [f_{20}^3]_3(\rho\ o\ )_3(a_3)\xi = [\mathrm{h\ o}^3]_3[\rho\mathrm{o}]_3(\mathrm{o}) \text{ and } \mathrm{b}_3 = \mathrm{b}_3 - g_3(P_3)$$

If $g_3(P_0) > b$, clearly, $(\rho_0)_3(a_3), 0)$ is unstable and the predator population is persistent. Assume $g_3(P_0)_3 < b$.

Then, $Q_3(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_3 + u_a = -(\mu_3(a_3) + (f_{30})_3)u_3 - (\alpha_3)_3(a_3)N_3$$

$$= -(\mu_3(a_3) + f_{30})u_3 - \alpha_3(a_3)\int_0^\infty u_3(t, a_3)da_3$$

$$u_3(t, 0) = \int_0^\infty \beta_3\,(a_3)u_3(t, a_3)da_3$$

Using the method developed by Cushing [2, 3], we suppose

$$u_3(t, a_3) = \tilde{u}_3(a_3)^{e^{\gamma(t-a_3)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_3}\,da_3$. Then,

A solution of the first differential equation for $\tilde{u}(a_3)$ gives,

$$\tilde{u}_3(a_3) = e^{-\int_0^a \mu_3(\tau_3)d\tau - (f_{30})_3 a}\left(\tilde{u}_3(0) - w\int_0^a e^{\int_0^\tau \mu_3(s_3)ds + (f_{30})}\tau_3\,(\alpha_3)_3(\tau_3)e^{\gamma_3\tau}d\tau\right).$$

Substituting the expression (13) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\right)\tilde{u}_3(0)$$

$$+w\int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\int_0^a (\alpha_3)_3(\tau_3)e^{-\gamma_3\tau_3}e^{(f_{30})_3\tau_3 + \int_0^\tau \mu_3(s_3)ds}d\tau_3\,da_3$$

$$= 0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(\tau_3)d\tau} da\right)(\tilde{u}_3(0))_3$$

$$+\left(-3-\int_0^\infty e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(\tau_3)d\tau}\int_0^a e^{(f_{30})_3\tau_3+\int_0^\tau \mu_3(s_3)ds+y\tau}\alpha_3(\tau)d\tau\, da\right)w=0. \quad (23)$$

Define $Y(\gamma,a)=e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(t_3)d\tau}$. Then

$$\left(1-\int_0^\infty \beta_3\,(a_3)Y_3(\gamma_3,a_3)da_3\right)u_3(0)$$

$$+w\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)\int_0^a(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\, da_3=0, \quad (24)$$

$$\left(\int_0^\infty Y_3(\gamma_3,a_3)da_3\right)\tilde{u}_3(0)$$

$$+\left(-1-\int_0^\infty Y_3(\gamma_3,a_3)\int_0^a(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\, da_3\right)w=0.$$

Equqtion (20) has a solution of the form of $u_3(t,a_3)=u_3(a_3)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_3(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_3(\gamma_3,a_3)=\int_0^a(\rho_0)_3(\tau_3)Y^{-3}\,(\gamma_3,a_3)d\tau_3$. We obtain the characteristic equation

$$f_{30_3}^l\left(\int_0^\infty Y_3(\gamma_3,a_3)da_3\right)\left(\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$+\left(1-\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)da\right)\left(1+(f_{30}^l)_3\int_0^\infty Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_3(a_3)=(\rho_0)_3(0)e^{-\int_0^a \mu_3(\tau_3)d\tau_3-(f_{30})_3a_3}$$

$$I_3(\gamma_3,a_3)=\int_0^a (\rho_0)_3(\tau_3)e^{\gamma_3\tau_3+(f_{30})_3\tau+\int_0^\tau \mu_3(s_3)ds_3}d\tau$$

$$=\frac{(\rho_0)_3(0)}{\gamma_3}(e^{\gamma_3a_3}-1)$$

Let

$$E_3(a_3)=e^{-(f_{30})_3a_3-\int_0^a \mu_3(\tau_3)d\tau_3}$$

Then $[Y_3]_3[\gamma_3]_3,a_3)=E_3(a_3)e^{-\gamma_3a_3}$. by a simple calculation,

$$\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)I(\gamma_3,a_3)da=\frac{[\rho_0]_3(0)}{\gamma}\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3a_3}\ da\right);$$

$$\int_0^\infty Y_3(\gamma_3,a_3)I(\gamma_3,a_3)da=\frac{[\rho_0]_3(0)}{\gamma_3}\left(\int_0^\infty E_3(a_3)e^{-\gamma_3a_3}\ da_3\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_3(0)}{\gamma_3}[f_{30}^3]\int_0^\infty E_3(a_3)e^{-\gamma_3a_3}da_3\ (\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3a_3}da_3)$$

$$+\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right).$$

$$\left(1+\frac{(\rho_0)_3(0)}{\gamma_3}(f_{30}^{'})_3\left(\int_0^\infty E_3(a_3)da_3-\int_0^\infty E_3(a_3)e^{-\gamma_3 a_3}da_3\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right)(1+\frac{(f_{30}^{'})_3}{\gamma_3}(\rho_0)_3(0)\int_0^\infty E_3(a_3)da_3)=0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_3(a_3)Y_3(\gamma_3 a_3)da_3 = 1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3 \neq 0,$$

Since $\int_0^\infty \beta_3(a_3)Y_3(\gamma_3 a_3)\int_0^a \alpha_3(\tau_3)Y^{-3}(\gamma_3,\tau_3)d\tau_3 da_3 > 0$. Therefore, the characteristic solution is

$$\gamma_3 = -(f_{30}^{'})_3(\rho_0)_3\int_0^\infty E_3(a_3)da_3.$$

Since $f_3'(P_0)_3 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_3), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_3), 0)$ which is related to the extinction of debit summation types

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_3 + (\mu_3(a_3) + [f_3]_3(P_3^*) + [f_2]_3(Q_3^*))P_3^* = 0, \quad (26)$$

$$p_3^*(0) = h_3(Q_3^*)\int_0^\infty \beta_3(a_3)p_3^*(a_J)p_3^*(a_3)da_3 \quad (27)$$

$$g_3(P_3^*) - b_3 = 0, \quad (28)$$

Where $P_3^* = \int_0^\infty p_3^*(a_3)\, da_3$

Equation (28) yields $P^* = g^{-3}(b)$,

From (26),

$$p_3^*(a) = p_3^*(0)e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}$$

Substituting into (27), we need to solve

$$\left(1 - h_3(Q_3^*)\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}da_3\right)p_3^*(0) = 0$$

For $Q_3^*$

Fix $p_3^*(0)$ and $P_3^*$ Let

$$H_3(z_3) = 1 - h_3(z_3)\int_0^{\infty} \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3} da_{,3}$$

And assume that

$$[k_3]_3 \int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(g_3^{-3}(b_3))a_3 - [k_2]_3 a_3}\, da_3 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_3(z_3)] > 0.$$

Clearly,

$$\frac{dH_3}{dz_3} = h_3(z)([f_2^{'}]_3(z_3))\int_0^{\infty} a_3\beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - f_3(P_3^*)a_3 - f_2(z_3)a_3} da_{\square}$$

$-h^{'}(z_3)\int_0^{\infty} \beta_3(a_3)\, e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3}\quad da_3$

By the assumptions on $\left[f_2^3\right]_3 and\ [h^3]^3 dH_3/dz > 0$. Hence, if

$H_3(0) = 1 - \int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3} da_3 < 0,$ (29)

There exists a unique positive solution $Q^*$

Now, we claim that $H_3(0) < 0$ follows if $g_3^{-}(b_3) < [f_3^{-3}]_3(r_3)$, where $r_3$ is the solution of (12).

Indeed, since

$$R_3(s_3) = \int_0^{\infty} \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 e^{-y\frac{a}{3}}} da_3$$

Is a decreasing function of $s_3$ when $P_3^* = [g^{-3}]_3(b_3) < [f_3^{-3}]_3(r_3)$, that is $[f_3]_3(P_3^*) < r$, $R\big(f_3(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_3(g^{-3}(b))) > 1, i.e.,$

$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3\left(g_3^{-3}(b)\right)a_3} > 1.$ (H4)

Then the system (4)-(6) has a unique positive equilibrium $(P_3^*(a_3), Q_3^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_3^*(a_3), Q_3^*)$ is more complicated. We again consider perturbations $m_3(t_3 a_3) = p_3(t_3, a_3) - p_3^*(a_3)$ and $W_3(t) = Q_3(t) - Q^*$.Then

$M_3(t) \equiv \int_0^{\infty} m_3(t, a_3) da_3 = P_3(t) - P_3^*,$

$$[m_3]_3 + [m_a]_3 = [p_3]_3 + [p_a]_3 - \left[p^*\right]_3(a_3) \approx -(\mu_3(a_3) + [f_{3*}]_3 + [f_{2*}]_3)m_3$$

$$-[f_{3*}^3]_3 p_3^*(a_3)M - [f_{2*}^3]_3 p_3^*(\mathrm{a})W_3,$$

$$m_3(t,0)p_3(t,0) - p_3^*(0) \approx h_* \int_0^{\infty} \beta_3(a_3) m_3(t, a_3)\, da + \frac{[h_*^3]_3}{[h_*]_3}\, p_3^*(0)W_3$$

$$\frac{dW_3}{dt} = \frac{dQ_3}{dt} \approx [g_*^3]_3 Q_3^* M_{,3}$$

Where $[f_{3*}]_3$ denoting $[f_3(P^*)]_3$, $[f_{3*}]_3$ denoting $[f_2(P^*)]_3$, $[f'_{3*}]_3$ denoting $[f'_3(P^*)]_3[f'_{2*}]_3$ denoting $[f_2]_3$ $(Q_3^*,\quad h_3^*$ denoting $h'_*$ denoting $h'_3$ $(Q_3^*)$, and $[g'_*]_3$

Denoting $[g']_3(P_3^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_3]_3 + [m_a]_3 = -(\mu_3(a_3) + [f_{3*}]_3 + [f_{2*}]_3)m_3 - [f'_{3*}]_3[p_3^*](a_3)M_3 - [f'_{2*}]_3 p_3^*(a_3)W_3$$

$$m_3(t,0) = [h_*]_3 \int_0^\infty \beta_3(a_3)m_3(t,a_3)da_3 + \frac{[h_*^3]_3}{[h^*]_3}[p^*]_3(0)\,W_{,3} \qquad (30)$$

$$\frac{dW_3}{dt} = [g_*^3]_3[Q_3^*]M_3$$

Integrating $[m_3]_3(t,a_3)$ with respect to a from 0 to ∞, we have

$$\frac{dM_3(t)}{dt} = \int_0^\infty [mt]_3(t,a_3)da_3$$

$$= -\int_0^\infty [m_a]_3\,(t,a_3)da_3 - \int_0^\infty \mu(a_3)m(t,a_3)da_3$$

$$-\left([f_{3*}]_3 + [f_{2*}]_3 + [f'_{3*}]_3\int_0^\infty p_3^*(a_3)da_3\right)M_3 - [f_{2*}^3]_3\int_0^\infty p_3^*(a_3)da_3 W_3$$

$$= m_3(t,0) - \int_0^\infty \mu_3(a_3)\,m_3(t,a_3)da_3$$

$$-\left([f_{3*}]_3 + [f_{2*}]_3 + [f'_{3*}]_3 P_3^*\right)M_3 - [f'_{2*}]_3 P_3^* W_{,3}$$

From (30), we arrive at the equations

$$\frac{dW_3(t)}{dt} = \int_0^{\infty} ([h_*]_3\beta_3(a_3) - \mu_3(a_3))m_3(t,a_3)da_3$$

$$-\left([f_{3*}] + [f_{2*}]_3 + \left[f_{3*}^{'}\right]_3 P_3^*\right) M_3 + \left(\frac{[h_*^{'}]_3}{[h_*]^3}\, P_3^*(0) - f_{2*}^{'} P_3^*\right) W_3,$$

(33)

$$\frac{dW_3(t)}{dt}\left[g^{'}{}_*\right]_3 Q_3^*\, M_{.3}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{n*}]_3\beta_3(a_3) - \mu_3(a_3))m_3(t,a_3)da_3 = ([h_*]_3\beta_3(\theta_3) - \mu_3(\theta_3))M_{,3}$$

Where $0< \theta_3 = \theta_3(t) < \infty$.

And if $a_3^2\,[a_3^2]_3<4\,[a_2]_3<4\,a_2 e$,

$\mathrm{Re}\,\lambda_3\left[\lambda_3\right]_3 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_3 = -\frac{3}{2}a_3[a2\ ]_3 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_*\,[h_*]_3\beta_3(a_3) - \mu_3(a_3)\right|$. Then$|n_3(t)|$ ⊠⊿∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_3^2 <4\,a_2 e$ and $[a_3] > 2\Delta$,or $[a_3^2]$ ⊠$[a_2]e$ and $(a_3 - \sqrt{(a_3^2)_3 - 4(a_2)_3 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+3}\right)$ and $\beta_3(a_3) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3) + d\tau_3 da} = \frac{3}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_3 + [\mu_0]$ 0.5. Assume $[f_3](x_3) = K_3 x_3$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_3(x_3) = [0x_3]/(3 + x_3)$; the functional responses $[[f_2](x) = [0]\, x_3/(1 + x_3)$ and $h_3(x_3)=(1+\omega_3 x_3)/(1 + x_3)$ (See Holling [7] and Rashevsky [13])

Equation (34),now has the form of

$$\frac{m_3 + [\mu_0] + [K][P_0])}{[m_3]_3 + (\mu_0)_3 + [K][P_0])_3^3\ [K]_3 [P_0]_3^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(3 + P_0) < b_3$ i e $[P_0] < b/(\delta_3 - b_3)$ then the equilibrium $(P_0)$ $[P_0]_3(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_3$

$[f_{3*}] = K_3\ P^*\ [f_{2*}] = \frac{\sigma Q_3^*}{3+Q_3^*} \quad [h_*] = \frac{3+\omega_3 Q_3^*}{3+Q_3^*}$

$[f_{3*}] = K_3\ [f_{2*}^3] = \frac{\sigma_3}{(3+Q_3^*)^2}\ [g_*^3] = \frac{\sigma_3}{(3+P_3^*)^2} \quad [h_*^3] = \frac{\omega_3 - 3}{(3+Q_3^*)^2}$

In order to have (H3) and (H4) hold, $m, \mu_0, K, b, \delta$, and $\omega$ need to satisfy

$$\frac{\omega_3}{(m+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma)^2}+\frac{2\omega_3}{([m_3]_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma_3)^3}<1$$

And

$$\frac{1}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^2}+\frac{2\omega_3}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}+(m_3+[\mu_0]_3+\frac{K_3b_3}{\delta_3-b_3}+\sigma_3)\,Q_3^*)^3$$

$$=m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}(1+Q^*)(1+\omega_3Q_3^*$$

Now

$$\left[a_3^2\right]-4\left[a_2\right]e=\left(\frac{2K_3b_3}{\delta_3-b_3}\right)^2+\frac{4K_3b_3\sigma_3\,Q_3^*}{(\delta_3-b_3)(1+Q_3^*)}+\frac{\sigma_3^2\,Q_3^{*2}}{(1+Q_3^*)^2}$$

$$\frac{4\,b_3\sigma_3\,(\delta_3-b)Q^*}{\delta_3(1+Q_3^*)^2}-\frac{4\,(\delta_3-b_3)^2(1-\omega_3)Q_3^*}{\delta_3(1+\omega_3Q_3^*)(1+Q_3^*)}P*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_3\,ae_3^{mu}\;\;[\mu_0]_3-\frac{a_3}{a_3+3}$$

Hence if $[a_3^2 > 4(a_2)_3 e$ and

$$[a_3] = \frac{2K_3 b_3}{\delta_3 - b_3} + \frac{\sigma_3 Q_3^*}{1 + Q_3^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_3]_3 = [f_{3*}]_3 + [f_{2*}]_3 + [f_{3*}] P_3^*$$

$$[a_2]_3 = \left[f_{2*}^3\right] P_3^* - \frac{h_*^3}{h_*} P^*(0)$$

$$e = [g_*^3]_3 Q_3^*$$

$$\eta_3(t) = \left[h_*\right]_3 \beta_3 \left(\theta_3(t)\right) - \mu_3(\theta(t)),$$

The system(33)can be written as the following system:

$$[M(t)]\left[a_3\right] M_3 - [a_2]_3 W_3 - [n(t)] M_3$$

$$w_3(t) = eM_3$$

Write

$$A = \begin{pmatrix} -[a_3]_3 - [a_2]_3 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_3 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_3 \\ W_3 \end{pmatrix}$$

The solution of (32)is

$$X_3(t) = e^{A,t}[X]_3(0) + \int_0 e^{At}$$

**This equation is useful in following type of problems . For actual model kindly see various papersThere is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant Λ, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive Λ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the Λ>0 generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible.**

**Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of (e) GW signals from BBHs with component spins $\chi_{1,2} \in [-0.99, 0.99]$ aligned (e&eb) with the orbital angular momentum, component masses $m_{1,2} \in [2,48] M_\odot$, and (e&eb) total mass $M_{total} \leq 50\ M_\odot$. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins >+0.6 and component masses $\gtrsim 20\ M_\odot$. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201 STOP Advanced LIGO may be the first experiment to (e) detect gravitational waves. Through superradiance of stellar black holes, it may also be (=) the first experiment to discover the QCD axion with (e&eb) decay constant above the GUT scale. When an axion's Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e) the black hole, forming a "gravitational atom." Through the superradiance process, the number of axions occupying (e) the bound levels grows exponentially, extracting (e) energy and angular momentum from the black hole. Axions**

**transitioning between levels of the gravitational atom and (e&eb) axions annihilating to (e) gravitons can produce (eb) observable gravitational wave signals. The signals are long-lasting, monochromatic, and can be distinguished from (e) ordinary astrophysical sources. Authors estimate up to O (1) transition events at aLIGO for (e) an axion between 10^-11 and (e&eb) 10^-10 eV and up to 10^4 annihilation events for (e) an axion between 10^-13 and 10^-11 eV. In the event of a null search, aLIGO can constrain (e) the axion mass for (e) a range of rapidly spinning black hole formation rates. Axion annihilations are also promising for (e) much lighter masses at future lower-frequency gravitational wave observatories; the rates have (e) large uncertainties, dominated by (e) supermassive black hole spin distributions. Projections for aLIGO are (=) robust against perturbations from (e) the black hole environment and account for (e) updated exclusion on (e&eb) the QCD axion of 6*10^-13 eV < ma < 2*10^-11 eV suggested by (e) stellar black hole spin measurements. Subjects: High Energy Physics - Phenomenology (hep-ph); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 91, 084011 (2015) DOI: 10.1103/PhysRevD.91.084011Cite as:arXiv:1411.2263 [hep-ph] (or arXiv:1411.2263v3 [hep-ph] for this version) Discovering the QCD Axion with Black Holes and Gravitational Waves Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang STOP Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric**

**charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by (e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave**

**signal. The signal sweeps upwards in frequency from 35 to 250 Hz with (e&eb) a peak gravitational-wave strain of (e) 1.0×10−21. It matches the waveform predicted by general relativity for (e) the inspiral and merger of a pair of black holes and the ringdown of the resulting (eb) single black hole. The signal was observed with a matched-filter signal-to-noise ratio of (e) 24 and a false alarm rate estimated to be (=) less than 1 event per 203 000 years, equivalent to (=) a significance greater than 5.1σ. The source lies at (eb) a luminosity distance of 410+160−180 Mpc corresponding to (e&eb) a redshift z=0.09+0.03−0.04. In the source frame, the initial black hole masses are (=) 36+5−4M⊙ and 29+4−4M⊙, and the final black hole mass is (=) 62+4−4M⊙, with (e&eb) 3.0+0.5−0.5M⊙c2 radiated in (eb) gravitational waves. All uncertainties define (eb) 90% credible intervals. These observations demonstrate (eb) the existence of binary stellar-mass black hole systems. Observation of Gravitational Waves from a Binary Black Hole Merger B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 061102 – Published 11 February 2016 This is the first direct detection of gravitational waves and (e&eb) the first observation of a binary black hole merger. Received 21 January 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.061102 STOP This is the second in a series of three papers on (e&eb) systematic analysis of rank 1 Coulomb branch geometries of (e) four dimensional N=2 SCFTs. In the first paper authors developed a strategy for (e) classifying physical rank-1 CB geometries of N=2 SCFTs. They show how to carry out this strategy computationally to construct (eb) the Seiberg-Witten curves and (e&eb) one-forms for all the rank-1 SCFTs. Explicit expressions are given for all cases, with the exception of the Nf=4 SU (2) gauge theory and (e&eb) the En SCFTs which were previously constructed. Classification includes all known rank-1 theories plus (+) a new one with (e&eb) an abelian flavor group, plus (+) nine additional theories whose existence is more speculative. Four of those, reported in first paper, depend on (=) the assumption of new frozen rank-1 SCFTs. Here authors also show that the assumption of the existence of certain rank-0 N=2 SCFTs leads to (eb) five additional consistent**

**rank-1 CB geometries. Subjects: High Energy Physics - Theory (hep-th) Cite as:arXiv: 1601.00011 [hep-th] (or arXiv: 1601.00011v2 [hep-th] for this version Geometric constraints on the space of N=2 SCFTs II: Construction of special Kähler geometries and RG flows Philip C. Argyres, Matteo Lotito, Yongchao Lü, Mario Martone (Submitted on 31 Dec 2015 (v1), last revised 6 Jan 2016 (this version, v2)) STOP Authors propose a model of Dark Energy in which the field currently dominating (e) the energy density of the universe is (=) an "axion field" linearly coupled to (e&eb) the Pontryagin density, tr (FΛF), (i.e., the exterior derivative of (e) the Chern-Simons form) of (e) a massive gauge field. Authors assume that the axion has (e) self-interactions corresponding to (e&eb) a non-trivial (exponential) potential. Authors argue that a non-vanishing magnetic helicity of (e) the gauge field triggers (eb) slow-rolling of the axion at field values far below (eb) the Planck scale. Proposal leads to (eb) a "Tracking Dark Energy Scenario" in which the contribution of the axion energy density to (e) the total energy density is (=) constant (and small) during (e&eb) the early radiation phase, until (e) a secular growth term proportional to (e) the Pontryagin density of (e) the gauge field becomes (=) dominant. The initially small contribution of (e) the axion field to (e) the total energy density is related to the observed small baryon-to-entropy ratio. Some speculations concerning the nature of the gauge field are offered.**

**Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Cite as:arXiv: 1601.00057 [hep-th] (or arXiv:1601.00057v2 [hep-th] for this version) Tracking Dark Energy from Axion-Gauge Field Couplings Stephon Alexander, Robert Brandenberger, Juerg Froehlich (Submitted on 1 Jan 2016 (v1), last revised 23 Jan 2016 (this version, v2)) STOP Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a**

**pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version) Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP In this work authors obtain bounds on (e&eb) the topological Abelian string-vortex and on (e&eb) the string-cigar, by using (e) a new measure of configurational complexity, known as (=) configurational entropy. In this way, the information-theoretical measure of (e) six-dimensional braneworld scenarios is capable to (e) probe situations where (e) the parameters responsible for (e) the brane thickness are arbitrary. The so-called configurational entropy (CE) selects (e&eb) the best value of the parameter in (eb) the model. This is accomplished by (e) minimizing the CE, namely, by (e) selecting the most appropriate parameters in the model that correspond to (e&eb) the most organized system, based upon (e) the Shannon information theory. This information-theoretical measure of complexity provides (eb) a complementary perspective to (e) situations where (e) strictly energy-based arguments are (=) inconclusive. Authors show that the higher the energy the higher the CE, what shows (eb) an important correlation between the energy of the localized field configuration and (e&eb) its**

**associated entropic measure. Subjects: High Energy Physics - Theory (hep-th) DOI:10.1016/j.physletb.2016.02.038 Cite as: arXiv: 1601.00076 [hep-th] (or arXiv: 1601.00076v2 [hep-th] for this version) Bounds on topological Abelian string-vortex and string-cigar from information-entropic measure R. A. C. Correa, D. M. Dantas, C. A. S. Almeida, Roldao da Rocha (Submitted on 1 Jan 2016 (v1), last revised 17 Feb 2016 (this version, v2)) STOP Sum of the Holevo quantity (that bounds the capacity of quantum channels to transmit (e&eb) classical information about an observable) and (e&eb) the quantum discord (a measure of the quantumness of (e) correlations of that observable) yields (eb) an observable-independent total given by (e) the quantum mutual information This split naturally delineates ,dilates and disseminates information about (e&eb) quantum systems accessible to (e) observers -- information that is redundantly transmitted by (e&eb) the environment -- while showing that it is (=) maximized for (e) the quasi-classical pointer observable. Other observables are accessible only via (e&eb) correlations with the pointer observable .Further authors prove an anti-symmetry property relating accessible information and (e&eb) discord. It shows that information becomes (=) objective -- accessible to (e) many observers -- only as quantum information is relegated to (e) correlations with (e&eb) the global environment, (eb) and, therefore, locally inaccessible. The resulting complementarity explains (eb) why, in a quantum Universe, we perceive objective classical reality while (e) flagrantly quantum superpositions are out of reach. Subjects: Quantum Physics (quant-ph) Journal reference: Scientific Reports 3, 1729 (2013) Report number: LAUR-13-26386 Cite as:arXiv: 1303.4659 [quant-ph] (or arXiv: 1303.4659v2 [quant-ph] for this version) Complementarity of quantum discord and classically accessible information Michael Zwolak, Wojciech H. Zurek STOP Niels Bohr proposed that the outcome of the measurement becomes (=) objective and real, and, hence, classical, when (e) its results can be communicated by (e) classical means. In this work authors revisit Bohr's postulate using (e) modern tools from (e) the quantum information theory. They find a**

**full confirmation of Bohr's idea: if a measurement device is in a nonclassical state, (eb) the measurement results cannot be communicated perfectly by classical means. In this case some part of information in the measurement apparatus is lost (e) in the process of communication: the amount of this lost information turns out to be (=) the quantum discord. The information loss occurs even when (e) the apparatus is not entangled with the system of interest. The tools presented in this work allow (eb) to generalize Bohr's postulate: authors show that for pure system-apparatus states (eb) quantum communication does not provide any advantage when (e) measurement results are communicated to (e) more than one recipient. They further demonstrate the superiority of quantum communication to (e&eb) two recipients on (e&eb) a mixed system-apparatus state and show that (eb) this effect is fundamentally different from (e) quantum state cloning. Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. Lett. 111, 040401 (2013) DOI:10.1103/PhysRevLett.111.040401 Report number: LA-UR-12-25351 Cite as: arXiv: 1304.2914 [quant-ph] (or arXiv: 1304.2914v2 [quant-ph] for this version) Quantum discord cannot be shared Alexander Streltsov, Wojciech H. Zurek STOP Wojciech H. Zurek, Don N. Page investigate thermodynamic equilibrium of (e) a self-gravitating perfect fluid in (eb) a spherically symmetric system containing (e) a black hole of mass M by means of (e) the Tolman-Oppenheimer-Volkoff (TOV) equation. At r >> 2M its solutions describe (eb) a black-body radiation atmosphere with (e&eb) the Hawking temperature T_BH~1/(8 \pi M) that is increasingly blueshifted as r approaches (eb) 2M. However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments:**

**Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by**

**(e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a firewall-like "Planck cocoon" with entropy close to (e) black hole entropy**

## Prescientiаl Dynamics And Contential Staticity Of Accentuation Attrition Models

## Module Four

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equailibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-445) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Some of the concepts to which model is applicable is as follows:**

**Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by (e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave signal. The signal sweeps upwards in frequency from 35 to 250 Hz with (e&eb) a peak gravitational-wave strain of (e) 1.0×10−21. It matches the waveform predicted by general relativity for (e) the inspiral and merger of a pair of black holes and the ringdown of the resulting (eb) single black hole. The signal was observed with a**

**matched-filter signal-to-noise ratio of (e) 24 and a false alarm rate estimated to be (=) less than 1 event per 203 000 years, equivalent to (=) a significance greater than 5.1σ. The source lies at (eb) a luminosity distance of 410+160−180 Mpc corresponding to (e&eb) a redshift z=0.09+0.03−0.04. In the source frame, the initial black hole masses are (=) 36+5−4M⊙ and 29+4−4M⊙, and the final black hole mass is (=) 62+4−4M⊙, with (e&eb) 3.0+0.5−0.5M⊙c2 radiated in (eb) gravitational waves. All uncertainties define (eb) 90% credible intervals. These observations demonstrate (eb) the existence of binary stellar-mass black hole systems. Observation of Gravitational Waves from a Binary Black Hole Merger B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 061102 – Published 11 February 2016 This is the first direct detection of gravitational waves and (e&eb) the first observation of a binary black hole merger. Received 21 January 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.061102 STOP This is the second in a series of three papers on (e&eb) systematic analysis of rank 1 Coulomb branch geometries of (e) four dimensional N=2 SCFTs. In the first paper authors developed a strategy for (e) classifying physical rank-1 CB geometries of N=2 SCFTs. They show how to carry out this strategy computationally to construct (eb) the Seiberg-Witten curves and (e&eb) one-forms for all the rank-1 SCFTs. Explicit expressions are given for all cases, with the exception of the Nf=4 SU (2) gauge theory and (e&eb) the En SCFTs which were previously constructed. Classification includes all known rank-1 theories plus (+) a new one with (e&eb) an abelian flavor group, plus (+) nine additional theories whose existence is more speculative. Four of those, reported in first paper, depend on (=) the assumption of new frozen rank-1 SCFTs. Here authors also show that the assumption of the existence of certain rank-0 N=2 SCFTs leads to (eb) five additional consistent rank-1 CB geometries. Subjects: High Energy Physics - Theory (hep-th) Cite as:arXiv: 1601.00011 [hep-th] (or arXiv: 1601.00011v2 [hep-th] for this version Geometric constraints on the space of N=2 SCFTs II: Construction of special Kähler geometries and RG flows Philip C. Argyres, Matteo Lotito, Yongchao Lü, Mario Martone (Submitted on 31 Dec 2015 (v1), last revised 6 Jan 2016 (this version, v2)) STOP Authors propose a model of Dark Energy in which the field currently dominating (e) the energy density of the universe is (=) an "axion field" linearly coupled to (e&eb) the Pontryagin density, tr (F∧F), (i.e., the exterior derivative of (e) the Chern-Simons form) of (e) a massive gauge field. Authors assume that the axion has (e) self-interactions corresponding to (e&eb) a non-trivial (exponential) potential. Authors argue that a non-vanishing magnetic helicity of (e) the gauge field triggers (eb) slow-rolling of the axion at field values far below (eb) the Planck scale. Proposal leads to (eb) a "Tracking Dark Energy Scenario" in which the contribution of the axion energy density to (e) the total energy density is (=) constant (and small) during (e&eb) the early radiation phase, until (e) a secular growth term proportional to (e) the Pontryagin density of (e) the gauge field becomes (=) dominant. The initially small contribution of (e) the axion field to (e) the total energy density is related to the observed small baryon-to-entropy ratio. Some speculations concerning the nature of the gauge field are offered.**

**Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Cite as:arXiv: 1601.00057 [hep-th] (or arXiv:1601.00057v2 [hep-th] for this version) Tracking Dark Energy from Axion-Gauge Field Couplings Stephon Alexander, Robert Brandenberger, Juerg Froehlich (Submitted on 1 Jan 2016 (v1), last revised 23 Jan 2016 (this version, v2)) STOP Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version) Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP In this work authors obtain bounds on (e&eb) the topological Abelian string-vortex and on (e&eb) the string-cigar, by using (e) a new**

**measure of configurational complexity, known as (=) configurational entropy. In this way, the information-theoretical measure of (e) six-dimensional braneworld scenarios is capable to (e) probe situations where (e) the parameters responsible for (e) the brane thickness are arbitrary. The so-called configurational entropy (CE) selects (e&eb) the best value of the parameter in (eb) the model. This is accomplished by (e) minimizing the CE, namely, by (e) selecting the most appropriate parameters in the model that correspond to (e&eb) the most organized system, based upon (e) the Shannon information theory. This information-theoretical measure of complexity provides (eb) a complementary perspective to (e) situations where (e) strictly energy-based arguments are (=) inconclusive. Authors show that the higher the energy the higher the CE, what shows (eb) an important correlation between the energy of the localized field configuration and (e&eb) its associated entropic measure. Subjects: High Energy Physics - Theory (hep-th) DOI:10.1016/j.physletb.2016.02.038 Cite as: arXiv: 1601.00076 [hep-th] (or arXiv: 1601.00076v2 [hep-th] for this version) Bounds on topological Abelian string-vortex and string-cigar from information-entropic measure R. A. C. Correa, D. M. Dantas, C. A. S. Almeida, Roldao da Rocha (Submitted on 1 Jan 2016 (v1), last revised 17 Feb 2016 (this version, v2)) STOP Sum of the Holevo quantity (that bounds the capacity of quantum channels to transmit (e&eb) classical information about an observable) and (e&eb) the quantum discord (a measure of the quantumness of (e) correlations of that observable) yields (eb) an observable-independent total given by (e) the quantum mutual information This split naturally delineates ,dilates and disseminates information about (e&eb) quantum systems accessible to (e) observers -- information that is redundantly transmitted by (e&eb) the environment -- while showing that it is (=) maximized for (e) the quasi-classical pointer observable. Other observables are accessible only via (e&eb) correlations with the pointer observable .Further authors prove an anti-symmetry property relating accessible information and (e&eb) discord. It shows that information becomes (=) objective -- accessible to (e) many observers -- only as quantum information is relegated to (e) correlations with (e&eb) the global environment, (eb) and, therefore, locally inaccessible. The resulting complementarity explains (eb) why, in a quantum Universe, we perceive objective classical reality while (e) flagrantly quantum superpositions are out of reach. Subjects: Quantum Physics (quant-ph) Journal reference: Scientific Reports 3, 1729 (2013) Report number: LAUR-13-26386 Cite as:arXiv: 1303.4659 [quant-ph] (or arXiv: 1303.4659v2 [quant-ph] for this version) Complementarity of quantum discord and classically accessible information Michael Zwolak, Wojciech H. Zurek STOP Niels Bohr proposed that the outcome of the measurement becomes (=) objective and real, and, hence, classical, when (e) its results can be communicated by (e) classical means. In this work authors revisit Bohr's postulate using (e) modern tools from (e) the quantum information theory. They find a full confirmation of Bohr's idea: if a measurement device is in a nonclassical state, (eb) the measurement results cannot be communicated perfectly by classical means. In this case some part of information in the measurement apparatus is lost (e) in the process of communication: the amount of this lost information turns out to be (=) the quantum discord. The information loss occurs even when (e) the apparatus is not entangled with the system of interest. The tools presented in this work allow (eb) to generalize Bohr's postulate: authors show that for pure system-apparatus states (eb) quantum communication does not provide any advantage when (e) measurement results are communicated to (e) more than one recipient. They further demonstrate the superiority of quantum communication to (e&eb) two recipients on (e&eb) a mixed system-apparatus state and show that (eb) this effect is fundamentally different from (e) quantum state cloning. Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. Lett. 111, 040401 (2013) DOI:10.1103/PhysRevLett.111.040401 Report number: LA-UR-12-25351 Cite as: arXiv: 1304.2914 [quant-ph] (or arXiv: 1304.2914v2 [quant-ph] for this version) Quantum discord cannot be shared Alexander Streltsov, Wojciech H. Zurek STOP Wojciech H. Zurek, Don N. Page investigate thermodynamic equilibrium of (e) a self-gravitating perfect fluid in (eb) a spherically symmetric system containing (e) a black hole of mass M by means of (e) the Tolman-Oppenheimer-Volkoff (TOV) equation. At r >> 2M its solutions describe (eb) a black-body radiation atmosphere with (e&eb) the Hawking temperature T_BH~1/(8 \pi M) that is increasingly blueshifted as r approaches (eb) 2M. However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments: Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by (e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a**

**firewall-like "Planck cocoon" with entropy close to (e) black hole entropy. Authors reproduce paper with an updated title and abstract. For a later study with similar results see G. 't Hooft, gr-qc/9706058 Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Physical Review D, Volume 29, Number 4, pp. 628 - 631 (1984) DOI:10.1103/PhysRevD.29.628 Cite as:arXiv:1511.07051 [gr-qc] (or arXiv:1511.07051v1 [gr-qc] for this version) Firewalls, black-hole thermodynamics, and singular solutions of the Tolman-Oppenheimer-Volkoff equation Wojciech H. Zurek, Don N. Page (Submitted on 22 Nov 2015) STOP Frederic Magniez, Ashwin Nayak consider the problem of testing the commutativity of (e) a black-box group specified by its k generators. The complexity (in terms of k) of (e) this problem was first considered by Pak, who gave (eb) a randomized algorithm involving O(k) group operations. They construct a quite optimal quantum algorithm for this problem whose complexity is in O~(k2/3)O~(k2/3). The algorithm uses and highlights the power of the quantization method of Szegedy. For the lower bound of Ω(k2/3)Ω(k2/3), authors give a reduction from (e) a special case of Element Distinctness to our problem. Along the way, they prove the optimality of the algorithm of (e) Pak for the randomized model. Algorithmica July 2007, Volume 48, Issue 3, pp 221-232 First online: 25 June 2007 Quantum Complexity of Testing Group Commutativity Frederic Magniez, Ashwin Nayak STOP Pawel Wocjan, Jon Yard analyze relationships between quantum computation and (e&eb) a family of generalizations of the Jones polynomial. Extending recent work by Aharonov et al., authors give efficient quantum circuits for (e) implementing the unitary Jones-Wenzl representations of the braid group. We use these to provide new quantum algorithms for (e) approximately evaluating a family of specializations of the HOMFLYPT two-variable polynomial of trace closures of braids. They also give (eb) algorithms for approximating the Jones polynomial of a general class of closures of (e) braids at roots of unity. Next authors provide (eb) a self-contained proof of a result of (e) Freedman et al. that any quantum computation can be replaced by (e&eb) an additive approximation of the Jones polynomial, evaluated at (eb) almost any primitive root of unity. Proof encodes two-qubit unitaries into (e&eb) the rectangular representation of (e) the eight-strand braid group. Authors then give QCMA-complete and (e&eb) PSPACE-complete problems which are based on braids. They conclude with direct proofs that evaluating the Jones polynomial of the plat closure at most primitive roots of unity is a #P-hard problem, while learning its most significant bit is PP-hard, circumventing (e) the usual route through the Tutte polynomial and graph coloring. Subjects: Quantum Physics (quant-ph); Geometric Topology (math.GT) Cite as: arXiv:quant-ph/0603069 (or arXiv:quant-ph/0603069v3 for this version) The Jones polynomial: quantum algorithms and applications in quantum complexity theory Pawel Wocjan, Jon Yard STOP Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan address the task of switching off the Hamiltonian of (e) a system by removing (e) all internal and system-environment couplings. Authors propose dynamical decoupling schemes, which use (e) only bounded-strength controls, for quantum many-body systems with local system Hamiltonians and local environmental couplings. To do so, they introduce (eb) the combinatorial concept of (e) balanced-cycle orthogonal arrays (BOAs) and show how to construct them from (e) classical error-correcting codes. The derived decoupling schemes may be useful as (=) a primitive for more complex schemes, e.g., for Hamiltonian simulation. For the case of n qubits and a 2-local Hamiltonian, the length of the resulting decoupling (e&eb) scheme scales as O (nlogn), improving over (e&eb) the previously best-known schemes that scaled quadratically with (e&eb) n. More generally, using balanced-cycle orthogonal arrays constructed from families of BCH codes, we show that bounded-strength decoupling for any ℓ-local Hamiltonian, where ℓ≥2, can be achieved using decoupling schemes of length at most O(nℓ−1logn). Subjects:Quantum Physics (quant-ph); Emerging Technologies (cs.ET) Journal reference: IEEE Trans. Inf. Theory, 62(5): 2881-2894, 2016 DOI: 10.1109/TIT.2016.2535183 Report number: MIT-CTP /4705 Cite as: arXiv: 1509.00408 [quant-ph] (or arXiv: 1509.00408v2 [quant-ph] for this version) Improved bounded-strength decoupling schemes for local Hamiltonians Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan STOP Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola propose dynamical control schemes for Hamiltonian simulation in many-body quantum systems that avoid instantaneous control operations and rely solely on (e) realistic bounded-strength control Hamiltonians. Each simulation protocol consists of (e) periodic repetitions of a basic control block, constructed as (=) a suitable modification of an "Eulerian decoupling cycle," that would otherwise implement (eb) a trivial (zero) target Hamiltonian. For an open quantum system coupled to (e) an uncontrollable environment, approach may be employed to (e) engineer an effective evolution that simulates (e&eb) a target Hamiltonian on the system, while suppressing (e) unwanted decoherence to the leading order. Authors present illustrative applications to (e) both closed- and open-system simulation settings, with (e&eb)**

**emphasis on simulation of non-local (two-body) Hamiltonians using (e) only local (one-body) controls. In particular, authors provide (eb) simulation schemes applicable to (e&eb) Heisenberg-coupled spin chains exposed to (e) general linear decoherence, and show how to simulate (e&eb) Kitaev's honeycomb lattice Hamiltonian starting from Ising-coupled qubits, as (=) potentially relevant to (e&eb) the dynamical generation of a topologically protected (e&eb) quantum memory. Additional implications for quantum information processing are discussed. Subjects: Quantum Physics (quant-ph) Journal reference: New J. Phys. 16, 045021 (2014) DOI: 10.1088/1367-2630/16/4/045021 Cite as:arXiv: 1310.4153 [quant-ph]**

**(or arXiv: 1310.4153v1 [quant-ph] for this version) Hamiltonian quantum simulation with bounded-strength controls Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola STOP Felix Fontein, Pawel Wocjan present a quantum algorithm for computing the period lattice of (e) infrastructures of fixed dimension. The algorithm applies to (e&eb) infrastructures that satisfy certain conditions. The latter are always fulfilled for (e) infrastructures obtained from (e) global fields, i.e., algebraic number fields and function fields with (e&eb) finite constant fields. The first main contributions is an exponentially better method for (e) sampling approximations of vectors of the dual lattice of (e) the period lattice than the methods outlined in the works of (e) Hallgren and Schmidt and Vollmer. This new method improves (eb) the success probability by (e) a factor of at least 2^ {n^2-1} where (e) n is the dimension. The second main contribution is a rigorous and complete proof that the running time of the algorithm is (=) polynomial in the logarithm of (e) the determinant of the period lattice and (e&eb) exponential in n. The third contribution is the determination of (e) an explicit lower bound on (e&eb) the success probability of our algorithm which greatly improves on (e&eb) the bounds given in the above works. The exponential scaling seems inevitable because (e) the best currently known methods for (e) carrying out fundamental arithmetic operations in infrastructures obtained from (e) algebraic number fields take exponential time. In contrast, the problem of computing (e&eb) the period lattice of infrastructures arising from function fields can be solved without (e) the exponential dependence on (e&eb) the dimension n since this problem reduces efficiently to the abelian hidden subgroup problem. This is also true for (e) other important computational problems in (eb) algebraic geometry. The running time of the best classical algorithms for (e) infrastructures arising from (e) global fields increases subexponentially with (e&eb) the determinant of the period lattice. Subjects: Quantum Physics (quant-ph) Cite as: arXiv: 1111.1348 [quant-ph]**

**(or arXiv: 1111.1348v2 [quant-ph] for this version) Quantum Algorithm for Computing the Period Lattice of an Infrastructure Felix Fontein, Pawel Wocjan STOP Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan advocate a new approach of (e) addressing hidden structure problems and (e&eb) finding efficient quantum algorithms. Authors introduce and investigate the Hidden Symmetry Subgroup Problem (HSSP), which is a generalization of the well-studied Hidden Subgroup Problem (HSP). Given a group acting on (e&eb) a set and an oracle whose level sets define (eb) a partition of the set; the task is to recover (e) the subgroup of symmetries of (e) this partition inside the group. The HSSP provides (eb) a unifying framework that, besides the HSP, encompasses (e) a wide range of algebraic oracle problems, including (e) quadratic hidden polynomial problems. While the HSSP can have (e) provably exponential quantum query complexity, authors obtain efficient quantum algorithms for various interesting cases. To achieve this, authors present a general method for (e) reducing the HSSP to the HSP, which works efficiently in several cases related to (e&eb) symmetries of polynomials. The HSSP therefore connects in a rather surprising way certain hidden polynomial problems with (e&eb) the HSP. Using this connection, authors obtain the first efficient quantum algorithm for (e) the hidden polynomial problem for (e) multivariate quadratic polynomials over fields of (e) constant characteristic. Authors also apply the new methods to (e&eb) polynomial function graph problems and (e&eb) present an efficient quantum procedure for (e) constant degree multivariate polynomials over (e&eb) any field. This result improves in several ways the currently known algorithms. Subjects: Quantum Physics (quant-ph) Journal reference: SIAM Journal on Computing 42:5 (2013) DOI: 10.1137/120864416 Cite as: arXiv: 1107.2189 [quant-ph]**

**(or arXiv: 1107.2189v2 [quant-ph] for this version) Hidden Symmetry Subgroup Problems Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan STOP Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan present an efficient general method for realizing a quantum walk operator corresponding to (e&eb) an arbitrary sparse classical random walk. Approach is based on Grover and Rudolph's method for preparing coherent versions of (e) efficiently integrable probability distributions. This method is intended for (e) use in quantum walk algorithms with polynomial speedups, whose complexity is usually measured in terms of (e&eb) how many times we have to apply a step of a quantum walk, compared to (e&eb) the number of necessary classical Markov chain steps. Authors consider a finer notion of complexity including (e) the**

**number of elementary gates it takes to implement each step of the quantum walk with (e&eb) some desired accuracy. The difference in complexity for (e) various implementation approaches is that method scales linearly in the sparsity parameter and (e&eb) poly-logarithmically with (e&eb) the inverse of the desired precision. Best previously known general methods either scale quadratically in (eb) the sparsity parameter, or polynomially in (eb) the inverse precision. Approach is especially relevant for (e) implementing quantum walks corresponding to (e&eb) classical random walks like those used in (eb) the classical algorithms for approximating permanents and sampling from binary contingency tables. In those algorithms, the sparsity parameter grows with (e&eb) the problem size, while (e) maintaining high precision is required. Subjects: Quantum Physics (quant-ph) Journal reference: Quantum Information and Computation, Vol.10, No.5&6, pp0420-0434 (2010) Cite as: arXiv: 0903.3465 [quant-ph] (or arXiv: 0903.3465v3 [quant-ph] for this version) Efficient Circuits for Quantum Walks Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan STOP Preparing the ground state of a system of interacting classical particles is (=) an NP-hard problem. Thus, there is in general no better algorithm to (e) solve this problem than exhaustively going through all N configurations of (e) the system to determine (eb) the one with lowest energy, requiring (e) a running time proportional to N. A quantum computer, if it could be built, could solve (e) this problem in time sqrt (N). Here, authors present a powerful extension of this result to (e) the case of interacting quantum particles, demonstrating (eb)that a quantum computer can prepare the ground state of (e) a quantum system as efficiently as it does for (e) classical systems. Subjects: Quantum Physics (quant-ph) Journal reference: Phys.Rev.Lett.102:130503, 2009 DOI:10.1103/PhysRevLett.102.130503 Cite as: arXiv: 0809.2705 [quant-ph] (or arXiv: 0809.2705v1 [quant-ph] for this version) Preparing ground states of quantum many-body systems on a quantum computer David Poulin, Pawel Wocjan STOP Daniel Nagaj, Pawel Wocjan construct a simple translationally invariant, nearest-neighbor Hamiltonian on (e&eb) a chain of 10-dimensional qudits that makes it possible to (e) realize universal quantum computing without (e) any external control during (e&eb) the computational process. Authors only require the ability to prepare an initial computational basis state which encodes (e) both the quantum circuit and its input. The computational process is (=) then carried out by (e) the autonomous Hamiltonian time evolution. After a time polynomially long in the size of (e) the quantum circuit has passed, the result of (e) the computation is obtained with (e&eb) high probability by (e) measuring a few qudits in the computational basis. This result also implies (eb) that there cannot (e) exist efficient classical simulation methods for (e) generic translationally invariant nearest-neighbor Hamiltonians on (e&eb) qudit chains, unless (e) quantum computers can be efficiently simulated by (e&eb) classical computers (or, put in complexity theoretic terms, unless BPP=BQP). Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. A 78, 032311 (2008) DOI: 10.1103/PhysRevA.78.032311 Cite as: arXiv: 0802.0886 [quant-ph] (or arXiv: 0802.0886v4 [quant-ph] for this version) Hamiltonian Quantum Cellular Automata in 1D Daniel Nagaj, Pawel Wocjan**

## 46. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_4}{dt} = -b_4 Q_4 + g_4(g_4)Q_4$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_4(0) = 0, \quad g_4(P_4) > 0, \; P_4 \; 0, \qquad \lim_{P_4 \to \infty} g_4(P_4) = (K_4)_4 \; \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_4 + (\rho_a)_4 + D_4(a_4, P_4, Q_4)\rho_4 + 0, \qquad 0 < a_4 \; \infty,$$

$$\rho_4(t, 0) = \int_0^\infty n_4(a_4, P_4, Q_4)\rho_4(t_4, a_4)da_4$$

Where $D_4$ is the mortality function and $n_4$ is the fertility function.

Put $_4$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_4}{dt}\right]_4 = x_4\left[\frac{dQ_2}{dt}\right]_4 - x_2\left[\frac{dQ_4}{dt}\right]_4$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_4 - x_4\left[\frac{dQ_2}{dt}\right]_4$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_4 - x_6\left[\frac{dQ_3}{dt}\right]_4$$

Where $x_4, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_4$ and $g_4(P_4)Q_4$

Now when you add the term $(b_4Q_4)$ with $g_4(P_4)(Q_4)$ we are infect see is that $(b_4Q_4)$ is credited incrementally part of $g_4(P_4)(Q_4)$ or instant nearly, when subtracted from $\frac{dQ_4}{dt}$, it means that $g_4(P_4)(Q_4)$ is slowly eaten by $\frac{dQ_4}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_4 = \mu_4(a_4) + (f_4)_4(P_4)_4 + (f_2)_4(Q)_4 \quad (2)$$

$$n_4 = \beta_4(a_4)h_4(Q_4), \quad (3)$$

Where $\mu_4$the age- dependent debit rate whit is $\mu_4(a_4)$ $\forall a > 0$; $(f_4)_4$ is the density- dependent death rate with $f_4(0) = 0$, $(f_4^4)_4(P_4)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_4)_4(P_4) = (k_4)_4$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_4(0) = 0, (f_2^4)_4(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_4(Q_4) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ $0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 4, h^4(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_4 + (\rho_a)_4 = -\mu_4(a_4) + (f_4)_4P_4) + \big(f_2)_4(Q_4)\big)(\rho_4)_4 \quad (4)$$

$$\rho_4(t_4, 0) = h_4(Q_4)\int_0^\infty \beta_4(a_4)\rho_4(t, a_4)da, \quad (5)$$

$$\frac{dQ_4}{dt}$$

$$= -b_4Q_4 + g_4(P_4)Q_4 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_4$ and $Q_4$ of (4-6) satisfy

$$0 \leq \rho_4(t, a_4) \leq \rho_4^*(t, a_4),$$

$$0 \leq Q_4(t) \leq Q_4^*(t),$$

Where $\rho_4^*(t, a_4)$ and $Q_4^*(t)$, are solutions of

$$(\rho_t^*)_4 + (\rho_a^*)_4 = -(\mu_4(a_4) + (f_4)_4(P_4^*) + (f_2)_4 Q_4^*))\, \rho_4^*, \quad (7)$$

$$\rho_4^*(t, 0) = \int_0^\infty \beta_4(a_4)\, \rho_4^*(t, a_4) da_4 \quad (8)$$

$$\frac{dQ_4^*}{dt} = -b_4\, Q_4^* + g_4(P_4^*)Q_4^*, \quad (9)$$

With $P_4^*(t) = \int_0^\infty \rho_4^*(t, a_4) da_4$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_4^*(t)}{dt} = P_4^*(t)\big(r_4 - (f_4)(P_4^*) - (f_2)_4(Q^*)\big), \quad (10)$$

$$\frac{dQ_4^*(t)}{dt} = Q_4^*(t)(-b_4 + g_4(P_4^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_4(a) e^{-\int_0^a \mu_4(\tau)dt} e^{-ra}\, da_4 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_4^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^{\infty} \beta_4(a_4) e^{-\int_0^a \mu_4(t)dt}\, da_4 > 1, \qquad (H1)$$

Which guarantees $r_4 > 0$.

*Remark.* If $\beta_4(a) = [\beta_0]_4 a_4 e^{-\alpha_4 a_4}, \mu_4(a_4) \equiv \mu_4$ then the solution of (12) is

$$r_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_4]_4(P_4) = 0, [f_2](Q_4) = \lambda_4 Q_4\ h_4(Q_4) = 1, \qquad and\ \ g_4(P_4) = c_4 P_4$$

$$[\rho_4]_4 + [\rho_a]_4 = -(\mu_4 + \lambda_4 Q_4)\rho_4$$

$$\rho_4(t,0) = \int_0^{\infty} [\beta_0]_4 a_4 e^{-\alpha_4 a_4} \rho_4(t, a_4)\, da_4$$

$$\frac{dQ_4}{dt} = -b_4 Q_4 + c_4 P_4 Q_4$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_4}{dt} = \omega_4 P_4 - \lambda_4 P_4 Q_4$$

$$\frac{dQ_4}{dt} = b_4 Q_4 - c_4 P_4 Q_4$$

Where $\omega_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$

## 47. THE EQUILIBRIUM $(\rho_0)_4(a_4), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_4(a_4), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_4 + (\rho_a)_4 = -\big(\mu_4(a_4) + (f_4)_4(P_4)\big)\rho,$$

$$\rho_4(t,0) = \int_0^\infty \beta_4(a_4)\rho_4(t,a_4)da_4 \qquad (13)$$

The solution of (13) has the form of

$$\rho_4(t,a_4) = Ce_4^{r(t-a_4)} - \int_0^a (\tau)d\tau - \int_0^4 (f_4)_4\big(P_4(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_4(t) = C_4 e^{r_4 t - \int_0^t (f_4)_4(P_4(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_4^{-4}(r)$ if $k_4 > r$. This is equivalent to

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{(k_4)_4 a_4} da_4 < 1.$$

In fact, if we define

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau)d\tau} e^{-s_4 a_4}\, da_4$$

$R(s)$ is a decreasing function. $R_4(r_4) = 1$ And $R_4(k_4)_4 < 1$ imply $r_4 < (k_4)_4$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_4(a_4), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{k_4 a_4} da_4 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_4 \equiv \int_0^\infty (\rho_0)_{44}(a_4)da_4$ *satisfies*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{-(f_4)_4(P_0)a_4} da = 1 \qquad (11)$$

Next, we study the stability of $(\rho_0)_4(a_4), 0)$

Let $u_4(t, a_4) = \rho(t, a_4) - \rho_0(a_4)$ Then

$$N_4(t) \equiv \int_0^\infty u_4(t, a_4)da_4 = \int_0^\infty \rho_4(t, a_4)da_4 - \int_0^\infty (\rho_0)_4(a_4)da = P_4(t) - P_0)_4$$

And

$$(u_4)_4 + (u_a)_4 = (\rho_t)_4 + (\rho_a)_4 - (\rho_0^4)_4(a_4)$$

$$= -\mu_4(a_4) + (f_4)_4(P_4) + (f_2)_4(Q_4))\,(u_4 + (\rho_0)_4) + \mu_4(a_4)(\rho_0)_4 + (f_4)_4(P_0)_4(\rho_0)_4$$

$$= -\mu_4(a_4)u_4 - (f_4)_4(P_4) + (f_2)_4(Q_4))\,u_4 - (\rho_0)_4\,(f_4)_4(P_4) - (f_4)_4(P_0)_4$$

$$+ (f_2)_4(Q_4)), \qquad (15)$$

Where $(\rho_0^4)_4(a_4)$ denotes $d(\rho_0)_4(a_4)/da_4$

Expansion $(f_4)_4(P_4)$ and $(f_2)_4(Q_4)$ at $(P_0)_4$ and 0 yield respectively

$$(f_4)_4(P_4) = (f_4)_4((P_0)_4 + N_4)$$

$$= (f_4)_4(P_0)_4 + (f_4^4)_4(P_0)_4 N_4 + o_4(N_4), \quad (16)$$

$$(f_2)_4(Q_4)$$

$$= (f_2^4)_4(0)Q + o_4(Q_4) \quad (17)$$

Write $f_4(P_0) = f_{40}, f_4'(P_0) = f_{40}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_4)_4 + (u_a)_4 \approx -(\mu_4(a_4) + (f_{40})_4)u - (f_{40}^4)_4(\rho_0)_4(N_4) - (f_{20}^4)_4(\rho_0)_4(Q_4)$$

$$u_4(t,0) = \rho_4(t,0)(\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a) + (\rho_0)_4(a_4)da_4 - (\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a_4)da_4 + (h_4(Q_4) - 4)(\rho_0)_4(0) \quad (18)$$

Expanding

$$h_4(Q_4) = 1 + (h^4)_4(0)Q_4 + 0(Q_4) \equiv 1 + h_0' Q_4 + 0Q_4$$

And substituting it into (48), we arrive at

$$u_4(t,0) \approx \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + (h_{a_4}')(\rho_0)_4 Q_4$$

The linearization is the following :

$$[u_4] + [u_4] = -(\mu_4(a_4) + [t_{40}]_4 u_4 - [\alpha_4]_4(a_4)N_4 - [\alpha_2]_4(a_4)Q_4$$

$$u_4(t,0) = \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + [\xi_4]Q_4 \quad (19)$$

$$\frac{dQ_4}{dt} = -b_4 Q_4$$

where

$$[\alpha_4]_4(a_4) = [t_4^4 a]_4[p, o]_4(a_4) d2(a_4) = [f_{20}^4]_4(\rho\, o\,)_4(a_4)\xi = [\mathrm{h\ o}^4]_4[\rho \mathrm{o}]_4(\mathrm{o}) \text{ and } \mathrm{b}_4 = \mathrm{b}_4 - g_4(P_4)$$

If $g_4(P_0) > b$, clearly, $(\rho_0)_4(a_4), 0)$ is unstable and the predator population is persistent. Assume $g_4(P_0)_4 < b$.

Then, $Q_4(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_4 + u_a = -(\mu_4(a_4) + (f_{40})_4)u_4 - (\alpha_4)_4(a_4)N_4$$

$$= -(\mu_4(a_4) + f_{40})u_4 - \alpha_4(a_4)\int_0^\infty u_4(t, a_4)da_4$$

$$u_4(t, 0) = \int_0^\infty \beta_4\,(a_4)u_4(t, a_4)da_4$$

Using the method developed by Cushing [2, 3], we suppose

$$u_4(t, a_4) = \tilde{\mathrm{u}}_4(a_4)^{e^{\gamma(t-a_4)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_4}\, da_4$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_4)$ gives,

$$\tilde{\mathrm{u}}_4(a_4) = e^{-\int_0^a \mu_4(\tau_4)d\tau - (f_{40})_4 a}\left(\tilde{\mathrm{u}}_4(0) - w\int_0^a e^{\int_0^\tau \mu_4(s_4)ds + (f_{40})}\tau_4\,(\alpha_4)_4(\tau_4)e^{\gamma_4\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1-\int_0^{\infty}\beta_4(a_4)e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\right)\tilde{u}_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a(\alpha_4)_4(\tau_4)e^{-\gamma_4\tau_4}e^{(f_{40})_4\tau_4+\int_0^\tau \mu_4(s_4)ds}d\tau_4\,da_4$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}da\right)(\tilde{u}_4(0))_4$$

$$+\left(-4-\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a e^{(f_{40})_4\tau_4+\int_0^\tau \mu_4(s_4)ds+y\tau}\alpha_4(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(t_4)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_4\,(a_4)Y_4(\gamma_4,a_4)da_4\right)u_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)\int_0^a(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4=0,$$

(24)

$$\left(\int_0^{\infty}Y_4(\gamma_4,a_4)da_4\right)\tilde{u}_4(0)$$

$$+\left(-1-\int_0^{\infty} Y_4(\gamma_4,a_4)\int_0^{a}(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4\right)w=0.$$

Equqtion (20) has a solution of the form of $u_4(t,a_4)=u_4(a_4)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_4(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_4(\gamma_4,a_4)=\int_0^a(\rho_0)_4(\tau_4)Y^{-4}(\gamma_4,a_4)d\tau_4$. We obtain the characteristic equation

$$f_{40}^{l}{}_4\left(\int_0^{\infty} Y_4(\gamma_4,a_4)da_4\right)\left(\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$+\left(1-\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)da\right)\left(1+(f_{40}^{l})_4\int_0^{\infty}Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_4(a_4)=(\rho_0)_4(0)e^{-\int_0^a\mu_4(\tau_4)d\tau_4-(f_{40})_4a_4}$$

$$I_4(\gamma_4,a_4)=\int_0^a(\rho_0)_4(\tau_4)e^{\gamma_4\tau_4+(f_{40})_4\tau+\int_0^{\tau}\mu_4(s_4)ds_4}d\tau$$

$$=\frac{(\rho_0)_4(0)}{\gamma_4}(e^{\gamma_4a_4}-1)$$

Let

$$E_4(a_4)=e^{-(f_{40})_4a_4-\int_0^a\mu_4(\tau_4)d\tau_4}$$

Then $[Y_4]_4[\gamma_4]_4, a_4) = E_4(a_4)e^{-\gamma_4 a_4}$. by a simple calculation,

$$\int_0^\infty \beta_4(a_4)Y_4(\gamma_4, a_4)I(\gamma_4, a_4)da = \frac{[\rho_0]_4(0)}{\gamma}\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}\ da\right);$$

$$\int_0^\infty Y_4(\gamma_4, a_4)I(\gamma_4, a_4)da = \frac{[\rho_0]_4(0)}{\gamma_4}\left(\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}\ da_4\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_4(0)}{\gamma_4}[f_{40}^4]\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}da_4\ (\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4)$$

$$+\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right).$$

$$\left(1 + \frac{(\rho_0)_4(0)}{\gamma_4}(f_{40}^{'})_4\left(\int_0^\infty E_4(\,a_4)da_4 - \int_0^\infty E_4(\,a_4)e^{-\gamma_4 a_4}da_4\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right)(1 + \frac{(f_{40}^{'})_4}{\gamma_4}(\rho_0)_4(0)\int_0^\infty E_4(\,a_4)da_4) = 0.$$

We need ũ(a) $\not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1 - \int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)da_4 = 1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4 \neq 0,$$

Since $\int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)\ \int_0^a \alpha_4\,(\tau_4)Y^{-4}(\gamma_4, \tau_4)d\tau_4 da_4 > 0$. Therefore, the characteristic solution is

$$\gamma_4 = -(f_{40}^{'})_4(\rho_0)_4 \int_0^{\infty} E_4(a_4)da_4.$$

Since $f_4'(P_0)_4 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_4), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_4), 0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_4 + (\mu_4(a_4) + [f_4]_4(P_4^*) + [f_2]_4(Q_4^*))P_4^* = 0, \quad (26)$$

$$p_4^*(0) = h_4(Q_4^*)\int_0^{\infty} \beta_4(a_4)p_4^*(a_J)p_4^*(a_4)da_4 \quad (27)$$

$$g_4(P_4^*) - b_4 = 0, \quad (28)$$

Where $P_4^* = \int_0^{\infty} p_4^*(a_4)\, da_4$

Equation (28) yields $P^* = g^{-4}(b)$,

From (26),

$p_4^*(a) = p_4^*(0)e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}$

Substituting into (27), we need to solve

$$\left(1 - h_4(Q_4^*)\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}da_4\right)p_4^*(0) = 0$$

For $Q_4^*$

Fix $p_4^*(0)$ and $P_4^*$ Let

$$H_4(z_4) = 1 - h_4(z_4)\int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}da_{,4}$$

And assume that

$$[k_3]_4\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(g_4^{-4}(b_4))a_4 - [k_2]_4 a_4}\, da_4 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_4(z_4)] > 0.$$

Clearly,

$$\frac{dH_4}{dz_4} = h_4(z)(\,[f_2']_4(z_4)\,)\int_0^\infty a_4\beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - f_4(P_4^*)a_4 - f_2(z_4)a_4}da_{\square}$$

$-h'(z_4)\int_0^\infty \beta_4(a_4)\, e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}\quad da_4$

the assumptions on $\left[f_2^4\right]_4 and\ [h^4]^4 dH_4/dz > 0$. Hence, if

$$H_4(0) = 1 - \int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4} da_4 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_4(0) < 0$ follows if $g_4^-(b_4) < [f_4^{-4}]_4(r_4)$, where $r_4$ is the solution of (12).

Indeed, since

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 e^{-y\frac{a}{4}} da_4}$$

Is a decreasing function of $s_4$ when $P_4^* = [g^{-4}]_4(b_4) < [f_4^{-4}]_4(r_4)$, that is $[f_4]_4(P_4^*) < r$, $R\big(f_4(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_4(g^{-4}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4\left(g_4^{-4}(b)\right)a_4} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_4^*(a_4), Q_4^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_4^*(a_4), Q_4^*)$ is more complicated. We again consider perturbations $m_4(t_4 a_4) = p_4(t_4, a_4) - p_4^*(a_4)$ and $W_4(t) = Q_4(t) - Q^*$.Then

$$M_4(t) \equiv \int_0^\infty m_4(t, a_4) da_4 = P_4(t) - P_4^*,$$

$$[m_4]_4 + [m_a]_4 = [p_4]_4 + [p_a]_4 - [p^*]_4(a_4) \approx -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4$$

$$-[f_{4*}^4]_4 p_4^*(a_4)M - [f_{2*}^4]_4 p_4^*(\mathrm{a})W_4,$$

$$m_4(t,0)p_4(t,0) - p_4^*(0) \approx h_* \int_0^\infty \beta_4(a_4)m_4(t,a_4)\,da + \frac{[h_*^4]_4}{[h_*]_4}\, p_4^*(0)W_4$$

$$\frac{dW_4}{dt} = \frac{dQ_4}{dt} \approx [g_*^4]_4 Q_4^* M_{,4}$$

Where $[f_{4*}]_4$ denoting $[f_4(P^*)]_4$, $[f_{4*}]_4$ denoting $[f_2(P^*)]_4$, $[f_{4*}^{'}]_4$ denoting $[f^{'}{}_4(P^*)]_4 [f_{2*}^{'}]_4$ denoting $[f_2]_4$ $(Q_4^*,$ $h_{*}$ denoting $h_*^{'}$ denoting $h_4^{'}$ $(Q_4^*)$, and $[g_*^{'}]_4$

Denoting $[g^{'}]_4(P_4^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_4]_4 + [m_a]_4 = -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4 - [f^{'}{}_{4*}]_4 [p_4^*](a_4)M_4 - [f_{2*}^{'}]_4 p_4^*(a_4)W_4$$

$$m_4(t,0) = [h_*]_4 \int_0^\infty \beta_4(a_4)m_4(t,a_4)da_4 + \frac{[h_*^4]_4}{[h^*]_4}[p^*]_4(0)\, W_{,4} \qquad (30)$$

$$\frac{dW_4}{dt} = [g_*^4]_4 [Q_4^*]M_4$$

This is equation is important in the following cases: problem of using observed correlations to infer causal relations is relevant to (e) a wide variety of scientific disciplines. Yet given correlations between just two classical variables, (eb) it is impossible to determine whether they arose from (e) a causal influence of one on the other or a common cause influencing both, unless one can implement (eb) a randomized intervention. **Katja Ried, Megan Agnew,**

**Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. Resch** here consider the problem of causal inference for (e) quantum variables. Authors introduce causal tomography, which unifies and generalizes (e&eb) conventional quantum tomography schemes to provide (eb) a complete solution to the causal inference problem using (e) a quantum analogue of a randomized trial. Authors furthermore show that, in contrast to the classical case, observed quantum correlations alone can sometimes provide (eb) a solution. They implement a quantum-optical experiment that allows (eb) to control (e&eb) the causal relation between two optical modes, and two measurement schemes -- one with (e&eb) and one without (e) randomization -- that extract this relation from (e) the observed correlations. Results show that entanglement and (e&eb) coherence, known to be central to (e) quantum information processing, also provide (eb) a quantum advantage for (e) causal inference. Subjects: Quantum Physics (quant-ph); Learning (cs.LG); General Relativity and Quantum Cosmology (gr-qc); Machine Learning (stat.ML) Journal reference: Nat Phys 11, 414-420 (2015) DOI: 10.1038/nphys3266 Cite as: arXiv: 1406.5036 [quant-ph] (or arXiv: 1406.5036v1 [quant-ph] for this version) **Inferring causal structure: a quantum advantage Katja Ried, Megan Agnew, Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. ReschSTOP D. T. Son, A. O. Starinets** review recent progress in applying the AdS/CFT correspondence to (e&eb) finite-temperature field theory. In particular, they show how the hydrodynamic behavior of field theory is reflected in (e&eb) the low-momentum limit of (e) correlation functions computed through (e&eb) a real-time AdS/CFT prescription, which they formulate. Authors also show how the hydrodynamic modes in field theory correspond to (e&eb) the low-lying quasinormal modes of (e) the AdS black p-brane metric. They provide a proof of the universality of (e) the viscosity/entropy ratio within (eb) a class of theories with (e&eb) gravity duals and formulate (eb) a viscosity bound conjecture. Possible implications for real systems are mentioned. Subjects: High Energy Physics - Theory (hep-th) Journal reference: Ann.Rev.Nucl.Part.Sci.57:95-118,2007 DOI: 10.1146/annurev.nucl.57.090506.123120Report number: INT PUB 07-02Cite as: arXiv:0704.0240 [hep-th]

(or arXiv:0704.0240v2 [hep-th] for this version) An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society **Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP** Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv:

1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version **Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP Viscosity, Black Holes, and Quantum Field Theory D. T. Son, A. O. Starinets STOP Authors** study information retrieval from evaporating black holes, assuming that the internal dynamics of a black hole is unitary and rapidly mixing, and assuming that the retriever has unlimited control over the emitted Hawking radiation. If the evaporation of the black hole has already proceeded past the ``half-way" point, where (e) half of the initial entropy has been radiated away, then (eb) additional quantum information deposited in (eb) the black hole is revealed in (eb) the Hawking radiation very rapidly. Information deposited prior to the half-way point remains concealed until (e) the half-way point, and then emerges (eb) quickly. These conclusions hold because (e) typical local quantum circuits are efficient encoders for (e) quantum error-correcting codes that nearly achieve (eb) the capacity of the quantum erasure channel. Estimate of a black hole's information retention time, based on (e) speculative dynamical assumptions, is just barely compatible with (e&eb) the black hole complementarity hypothesis. **Black holes as mirrors: quantum information in random subsystems Patrick Hayden1 and John Preskill2 Published 26 September 2007 Journal of High Energy Physics, Volume 2007, JHEP09 (2007) STOP** By analyzing the key properties of black holes from the point of (e) view of quantum information, authors derive a model-independent picture of (e) black hole quantum computing. It has been noticed that this picture exhibits (eb) striking similarities with (e&eb) quantum critical condensates, allowing (eb) the use of a common language to describe (eb) quantum computing in both systems. Authors analyze such quantum computing by allowing coupling to (e&eb) external modes, under the condition that the external influence must be (=) soft-enough in order not to (e) offset the basic properties of the system. Authors derive model-independent bounds on some crucial time-scales, such as (=) the times of gate operation, decoherence, maximal entanglement and total scrambling. They show that for black hole type quantum computers all these time-scales are of (e) the order of the

black hole half-life time. Furthermore, authors construct explicitly a set of Hamiltonians that generates (eb) a universal set of quantum gates for (e) the black hole type computer. They find that the gates work at maximal energy efficiency. Furthermore, authors establish a fundamental bound on the complexity of quantum circuits encoded on (eb&e) these systems, and characterize the unitary operations that are (=) implementable. It becomes apparent that the computational power is very limited due to (e) the fact that the black hole life-time is of the same order of (e) the gate operation time. As a consequence, it is impossible to retrieve its information, within (eb) the life-time of a black hole, by externally coupling to (e&eb) the black hole qubits. However, authors show that, in principle, coupling to (e&eb) some of the internal degrees of freedom allows (eb) acquiring knowledge about the micro-state. Still, due to the trivial complexity of operations that can be performed, (eb) there is no time advantage over the collection of Hawking radiation and subsequent decoding. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Report number: LMU-ASC 03/16, MPP-2016-1 Cite as: arXiv:1605.01407 [hep-th] (or arXiv:1605.01407v1 [hep-th] for this version) **Universality of Black Hole Quantum Computing Gia Dvali, Cesar Gomez, Dieter Lust, Yasser Omar, and Benedikt Richter (Submitted on 4 May 2016) STOP** In this addendum to [arXiv: 1402.5674] two points are discussed. In the first additional evidence is provided (eb) for a dual connection between the geometric length of an Einstein-Rosen bridge and (e&eb) the computational complexity of the quantum state of the dual CFT's. The relation between growth of complexity and (e&eb) Page's “Extreme Cosmic Censorship” principle is also remarked on. The second point involves a gedanken experiment in which Alice measures (eb) a complete set of commuting observables at her end of (e) an Einstein-Rosen bridge is discussed. An apparent paradox is resolved by appealing to (e) the properties of GHZ tripartite entanglement. **Addendum to computational complexity and black hole horizons Leonard Susskind* Version of Record online: 18 JAN 2016 DOI: 10.1002/prop.201500093 STOP** Analyzing some well established facts, authors give a model-independent

parameterization of black hole quantum computing in terms of (e&*eb) a set of macro and micro quantities and their relations. These include the relations between the extraordinarily-small energy gap of (e) black hole qubits and important time-scales of (e) information-processing, such as, scrambling time and Page's time. Authors then show, confirming and extending previous results, that other systems of nature with (e&eb) identical quantum informatics features are (=) attractive Bose-Einstein systems at (eb) the critical point of quantum phase transition. Here authors establish a complete isomorphy between (e&eb) the quantum computational properties of these two systems. In particular, they show that the quantum hair of a critical condensate is (=) strikingly similar to (=) the quantum hair of a black hole. Irrespectively whether one takes the similarity between the two systems as a remarkable coincidence or as (=) a sign of a deeper underlying connection, the following is evident. Black holes are not unique in their way of (e) quantum information processing and they can manufacture black hole based quantum computers in labs by (e) taking advantage of quantum criticality. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Atomic Physics (physics. Atom-ph); Quantum Physics (quant-ph) Cite as: arXiv: 1601.01329 [hep-th] (or arXiv: 1601.01329v1 [hep-th] for this version) **Black Hole Based Quantum Computing in Labs and in the Sky Gia Dvali, Mischa Panchenko (Submitted on 6 Jan 2016) STOP Gia Dvali, Alexander Gußmann** show that the existence of black holes with (e&eb) classical skyrmion hair invalidates (e) standard proofs that global charges, such as (=) the baryon number, cannot be conserved by (e) a black hole. By carefully analyzing the standard arguments based on (e) a Gedankenexperiment in which a black hole is seemingly-unable to (e) return the baryon number that it swallowed, authors identify inconsistencies in this reasoning, which does not (e) take into the account neither the existence of skyrmion black holes nor the baryon/skyrmion correspondence. Authors then perform a refined Gedankenexperiment by incorporating (e) the new knowledge and show that (eb) no contradiction with conservation of baryon number takes place at (eb) any stage of black hole evolution. Analysis also indicates (eb) no conflict

between semi-classical black holes and (e&eb) the existence of baryonic gauge interaction arbitrarily-weaker than gravity. Next, authors study classical cross sections of a minimally-coupled massless probe scalar field scattered by (e) _ a skyrmion black hole. Authors investigate how the skyrmion hair manifests itself by (e) comparing this cross section with (e&eb) the analogous cross section caused by (e) a Schwarzschild black hole which has (e) the same ADM mass as (=) the skyrmion black hole. Here authors find an order-one difference in the positions of (e) the characteristic peaks in (eb) the cross sections. The peaks are shifted to smaller scattering angles when (e) the skyrmion hair is present. This comes from the fact that the skyrmion hair changes (e&eb) the near horizon geometry of the black hole when compared to (e&eb) a Schwarzschild black hole with same ADM mass. Authors keep the study of this second aspect general so that (e) the qualitative results obtained can also be applied to(e&eb) black holes with classical hair of different kind. Subjects: High Energy Physics - Theory (hep-th); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU-ASC 18/16 Cite as: arXiv:1605.00543 [hep-th] (or arXiv:1605.00543v2 [hep-th] for this version) A preliminary discussion is given of (e) the prospects that gravitational-wave observations of (e) binary inspiral of black holes could reveal or constrain (e&eb) quantum modifications to (e) black hole dynamics, such as are required to (e) preserve postulates of quantum mechanics. Different proposals for such modifications are characterized by (e) different scales, and the size of these scales relative to (e&eb) those probed by observation of (e) inspiral signals is important in determining (eb) the feasibility of finding experimental signatures. Certain scenarios with strong quantum modifications in (eb) a region extending well outside the horizon are expected to modify (e&eb) classical evolution, and distort (e&eb) the near-peak gravitational wave signal, suggesting (eb) a search for departures from (e) waveforms predicted by general relativity. The near agreement of the GW150914 signals with (e&eb) such waveforms is discussed, and indicates (eb) constraints on some such scenarios. Important strategies for more precise future tests are 1) to develop more precise

predictions from (e) scenarios proposing quantum modifications, and 2) searching for observed deviations from (e) numerical relativity predictions via (e&eb) analysis of gravity wave data, particularly focussing on (e&eb) the signal region corresponding to plunge and merger. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1602.03622 [gr-qc] (or arXiv: 1602.03622v2 [gr-qc] for this version) **Gravitational wave tests of quantum modifications to black hole structure -- with post-GW150914 update Steven B. Giddings STOP** Direct LHC bounds on (e&eb) colored SUSY particles now corner naturalness more than the measured value of the Higgs mass does. Bounds on the gluino are of (e) particular importance, since it radiatively "sucks" up (e) the stop and Higgs-up soft masses. As a result, even models that easily accommodate a 125 GeV Higgs are almost as tuned as (=) the simplest version of SUSY, the MSSM: at best at the percent level. In this paper, authors further examine how current LHC results constrain (e) naturalness in three classes of models that may relax LHC bounds on (e&eb) sparticles: split families (e&eb), baryonic RPV (e&eb), and Dirac gauginos. In models of split families and bRPV, the bounds on the gluino are (=) only slightly reduced, resulting in (eb) a few percent tuning. In particular, having a natural spectrum in bRPV models typically implies (eb) that tops, Ws, and Zs are easily produced (eb) in the cascade decays of (e) squarks and gluinos. The resulting leptons and missing energy push (eb) the gluino mass limit above 1 TeV. Even when the gluino has a Dirac mass and does not contribute to the stop mass at one loop, tuning reappears in calculable models because there is no symmetry imposing the supersoft limit. Authors conclude that, even if sparticles are found at LHC-14, naturalness will not emerge triumphant. Subjects: High Energy Physics - Phenomenology (hep-ph)DOI:10.1007/JHEP03(2014)022 Cite as:rXiv:1309.3568 [hep-ph] (or arXiv:1309.3568v2 [hep-ph] for this version) **The Last Vestiges of Naturalness Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang, Ken Van Tilburg, Giovanni Villadoro STOP** The lack of evidence for new physics beyond (e) the standard model at the LHC points to (e) a paucity of new particles near the weak scale. This suggests that the weak scale is tuned and that supersymmetry, if present at

all, is realized (eb) at higher energies. The measured Higgs mass constrains (e) the scalar sparticles to be below 10^5 TeV, while (e) gauge coupling unification favors (eb) Higgsinos below 100 TeV. Nevertheless, in many models gaugino masses are (=) suppressed and remain within reach of the LHC. Tuning the weak scale and (e&eb) the renormalization group evolution of (e) the scalar masses constrain (e) Split model building. Due to the small gaugino masses, either the squarks or the up-higgs often run tachyonic; in the latter case, successful electroweak breaking requires heavy higgsinos near the scalar sparticles. We discuss the consequences of tuning the weak scale and the phenomenology of several models of Split supersymmetry including anomaly mediation, U(1)_(B-L) mediation, and Split gauge mediation. Subjects: High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Experiment (hep-ex) DOI: 10.1007/JHEP02 (2013)126 Cite as: arXiv: 1210.0555 [hep-ph] (or arXiv: 1210.0555v2 [hep-ph] for this version) **Mini-Split Asimina Arvanitaki, Nathaniel Craig, Savas Dimopoulos, Giovanni Villadoro STOP** Light states associated with (e&eb) the hierarchy problem affect (e&eb) the Higgs LHC production and decays. Authors illustrate this within the MSSM and two simple extensions applying (e&eb) the latest bounds from (e) LHC Higgs searches. Large deviations in the Higgs properties are expected in (eb) a natural SUSY spectrum. The discovery of a non-Standard-Model Higgs may signal (eb) the presence of light stops accessible at (eb) the LHC. Conversely, the more the Higgs is Standard-Model-like, (eb) the more tuned the theory becomes. Taking the ratio of different Higgs decay channels at the LHC cancels (er&eb) the leading (eb) QCD uncertainties and potentially improves (eb) the accuracy in Higgs coupling measurements to (e) the percent level. This may lead to (eb) the possibility of doing precision Higgs physics at (eb) the LHC. Finally, authors entertain the possibility that the ATLAS excess around (e&eb) 125 GeV persists with (e&eb) a Higgs production cross-section that is enhanced compared to (e&eb) the SM. This increase can only be accommodated in extensions of (e) the MSSM and it may suggest (eb) that stops lie below 400 GeV, likely within (eb) reach of next year's LHC run. Subjects: High Energy Physics - Phenomenology (hep-ph) DOI: 10.1007/JHEP02 (2012)144 Cite

as: arXiv: 1112.4835 [hep-ph] (or arXiv: 1112.4835v2 [hep-ph] for this version) **A Non Standard Model Higgs at the LHC as a Sign of Naturalness Asimina Arvanitaki, Giovanni Villadoro STOP** It has recently been suggested that the presence of a plenitude of light axions, an Axiverse, is (=) evidence for the extra dimensions of string theory. Authors discuss the observational consequences of (e) these axions on astrophysical black holes through (e&eb) the Penrose superradiance process. When an axion Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e&eb) the black hole "nucleus" forming (eb) a gravitational atom in the sky. The occupation number of superradiant atomic levels fed by (e) the energy and angular momentum of the black hole grows exponentially. The black hole spins down and an axion Bose-Einstein condensate cloud forms (eb) around it. When the attractive axion self-interactions become (=) stronger than the gravitational binding energy, (eb) the axion cloud collapses, a phenomenon known in condensed matter physics as "Bosenova". The existence of axions is first diagnosed by (e) gaps in the mass vs (e&eb) spin plot of astrophysical black holes. For young black holes the allowed values of (e) spin are quantized, giving rise to (eb) "Regge trajectories" inside the gap region. The axion cloud can also be observed directly either through (e&eb) precision mapping of the near horizon geometry or through (e&eb) gravitational waves coming from (e) the Bosenova explosion, as well as axion transitions and (e&eb) annihilations in (eb) the gravitational atom. Estimates suggest (eb) that these signals are detectable in (eb) upcoming experiments, such as Advanced LIGO, AGIS, and LISA. Current black hole spin measurements imply (eb) an upper bound on (e&eb) the QCD axion decay constant of (e) 2 x 10^17 GeV, while (e) Advanced LIGO can detect (e&eb) signals from (e) a QCD axion cloud with (e&eb) a decay constant as low as (=) the GUT scale. Authors finally discuss the possibility of observing the gamma-rays associated with the Bosenova explosion and, perhaps, the radio waves from axion-to-photon conversion for the QCD axion. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Journal

reference:Phys.Rev.D83:044026, 2011 DOI: 10.1103/PhysRevD.83.044026 Cite as: arXiv: 1004.3558 [hep-th] (or arXiv: 1004.3558v2 [hep-th] for this version) **Exploring the String Axiverse with Precision Black Hole Physics Asimina Arvanitaki, Sergei Dubovsky (Submitted on 20 Apr 2010 (v1), last revised 16 Aug 2011 (this version, v2)) STOP Skyrmion Black Hole Hair: Conservation of Baryon Number by Black Holes and Observable Manifestations Gia Dvali, Alexander Gußmann STOP Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust** discuss BMS supertranslations both at (eb) null-infinity and on the horizon for the case of the Schwarzschild black hole. Authors show that both kinds of supertranslations lead to (eb) infinitely many gapless physical excitations. On this basis authors construct a quotient algebra using (e) suited superpositions of (e) both kinds of transformations which cannot be compensated by (e) an ordinary BMS-supertranslation and therefore are (=) intrinsically due to (e) the presence of an event horizon. Authors show that these quotient transformations are (=) physical and generate (eb) gapless excitations on the horizon that can account for (e) the gravitational hair as well as for (e) the black hole entropy. They identify the physics of these modes as associated with (e&eb) Bogolioubov-Goldstone modes due to quantum criticality. Classically the number of these gapless modes is (=) infinite. However, authors show that due to (e) quantum criticality the actual amount of information-carriers becomes (=) finite and consistent with Bekenstein entropy. Although they only consider the case of Schwarzschild geometry, the arguments are extendable to(e&eb) arbitrary space-times containing event horizons. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU--ASC 05/16, MPP--2016—3 Cite as: arXiv: 1601.03725 [hep-th] (or arXiv: 1601.03725v1 [hep-th] for this version) **Gravitational Black Hole Hair from Event Horizon Supertranslations Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust STOP Gia Dvali, Cesar Gomez, Nico Wintergerst** suggest that holography can be formulated in terms of (e&eb) the information capacity of (e) the St\"uckelberg degrees of freedom that maintain

gauge invariance of (e) the theory in the presence of (e) an information boundary. These St\"uckelbergs act as (e&eb) qubits that account for a certain fraction of (e) quantum information. Their information capacity is measured by (e) the ratio of the inverse St\"uckelberg energy gap to (e) the size of the system. Systems with (e&eb) the smallest gap are (=) maximally holographic. For massless gauge systems this information measure is (=) universally equal to (=) the inverse coupling evaluated at the systems' length scale. In this language it becomes very transparent why the St\"uckelberg information capacity of black holes saturates (e full) the Bekenstein bound and accounts for (e) the entire information of the system. The physical reason is that the strength of quantum interaction is bounded from below by the gravitational coupling, which scales as (=) area. Observing the striking similarity between the scalings of the energy gap of the boundary St\"uckelberg modes and (e&eb): (AC=DC=0) the Bogoliubov modes of critical many-body systems, authors establish a connection between holography and (e&eb) quantum criticality through (e&eb) the correspondence between these modes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1511.03525 [hep-th] (or arXiv: 1511.03525v1 [hep-th] for this version) **Stückelberg Formulation of Holography Gia Dvali, Cesar Gomez, Nico Wintergerst STOP B**lack hole entropy, denoted by N, in (semi)classical limit is (=) infinite. This scaling reveals (eb) a very important information about the qubit degrees of (e) freedom that carry black hole entropy. Namely, the multiplicity of qubits scales as N, whereas (e) their energy gap and their coupling as 1/N. Such a behavior is indeed exhibited by (e) Bogoliubov-Goldstone degrees of freedom of (e) a quantum-critical state of N soft gravitons (a condensate or a coherent state) describing (eb) the black hole quantum portrait. They can be viewed as (=) the Goldstone modes of (e) a broken symmetry acting on (e&eb) the graviton condensate. In this picture Minkowski space naturally emerges as (=) a coherent state of infinite-N gravitons of (e) infinite wavelength and it carries (e&eb) an infinite entropy. In this paper authors **Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015)** ask what is the geometric meaning (if any) of (e) the classical limit of this symmetry. Authors argue that the infinite-

N limit of Bogoliubov-Goldstone modes of critical graviton condensate is described by (e) recently-discussed classical BMS super-translations broken by the black hole geometry. However, the full black hole information can only be recovered for (e) finite N, since the recovery time becomes (=) infinite in classical limit in which N is (=) infinite. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) DOI: 10.1016/j.physletb.2015.11.073 Report number: MPP-2015-215, LMU-ASC 59/15 Cite as: arXiv: 1509.02114 [hep-th] (or arXiv: 1509.02114v1 [hep-th] for this version) **Classical Limit of Black Hole Quantum N-Portrait and BMS Symmetry Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015) STOP** Recent ideas about understanding physics of black hole information-processing in terms of (e&eb) quantum criticality allow us to (eb) implement black hole mechanisms of quantum computing within (eb) critical Bose-Einstein systems. The generic feature, uncovered both by analytic and numeric studies, is (=) the emergence at the critical point of (e) gapless weakly-interacting modes, which act as (=) qubits for information-storage at a very low energy cost. These modes can be effectively described in terms of (e&eb) either Bogoliubov or Goldstone degrees of freedom. The ground-state at the critical point is (=) maximally entangled and far from being classical. We confirm this near-critical behavior by a new analytic method. **Gia Dvali, Mischa Panchenko** compute growth of entanglement and show its consistency with (e&eb) black hole type behavior. On the other hand, in the over-critical regime the system develops (eb) a Lyapunov exponent and scrambles quantum information very fast. By, manipulating (e&eb) the system parameters externally, authors can put it in and out of various regimes and in this way control (e&eb) the sequence of information storage and processing. By using gapless Bogoliubov modes as (=) control qubits they design some simple logic gates. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Cite as: arXiv: 1507.08952 [hep-th] (or arXiv: 1507.08952v1 [hep-th] for this version) **Black Hole Type Quantum**

**Computing in Critical Bose-Einstein Systems Gia Dvali, Mischa Panchenko (Submitted on 31 Jul 2015) STOP Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst** establish a fundamental connection between quantum criticality of a many-body system, such as Bose-Einstein condensates, and (e&eb) its capacity of information-storage and processing. For deriving the effective theory of modes in the vicinity of the quantum critical point they develop a new method by mapping a Bose-Einstein condensate of N-particles onto (e&eb) a sigma model with a continuous global (pseudo)symmetry that mixes (e&eb) bosons of different momenta. The Bogolyubov modes of the condensate are mapped onto (e&eb) the Goldstone modes of the sigma model, which become (=) gapless at the critical point. These gapless Goldstone modes are (=) the quantum carriers of information and entropy. Analyzing their effective theory, authors observe the information-processing properties strikingly similar to (=) the ones predicted by the black hole portrait. The energy cost per qubit of information-storage vanishes in (eb) the large-N limit and the total information-storage capacity increases with (e&eb) N either exponentially or as a power law. The longevity of information-storage also increases with (e&eb) N, whereas (e) the scrambling time in (eb) the over-critical regime is controlled by (e&eb) the Lyapunov exponent and scales logarithmically with (e&eb) N. This connection reveals (eb) that the origin of black hole information storage lies in (eb) the quantum criticality of the graviton Bose-gas, and that much simpler systems that can be manufactured in table-top experiments can exhibit (eb) very similar information-processing dynamics. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 92, 125002 (2015) DOI: 10.1103/PhysRevD.92.125002 Cite as: arXiv:1507.02948 [hep-th] (or arXiv:1507.02948v1 [hep-th] for this version) **Nambu-Goldstone Effective Theory of Information at Quantum Criticality Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst STOP Gia Dvali, Cesar Gomez** show that a positive cosmological constant is (=) incompatible with the quantum-corpuscular resolution of de Sitter metric in form of (e) a coherent state. The reason is very general and is due to (e) the

quantum self-destruction of (e) the coherent state because of (e) the scattering of constituent graviton quanta. This process creates an irreversible quantum clock, which precludes (e) eternal de Sitter. It also eliminates (e) the possibility of Boltzmann brains and Poincare recurrences. This effect is expected to be (=) part of any microscopic theory that takes into account (e) the quantum corpuscular structure of the cosmological background. This observation puts the cosmological constant problem in a very different light, promoting it, from a naturalness problem, into (e&eb) a question of quantum consistency. We are learning that quantum gravity cannot tolerate (e) exceedingly-classical sources. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1412.8077 [hep-th]

(or arXiv:1412.8077v1 [hep-th] for this version) **Quantum Exclusion of Positive Cosmological Constant? Gia Dvali, Cesar Gomez (Submitted on 27 Dec 2014) STOP G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger** establish a connection between the ultra-Planckian scattering amplitudes in field and string theory and (e&eb) unitarization by black hole formation in these scattering processes. **Using as a guideline an explicit microscopic theory in which the black hole represents (eb) a bound-state of many soft gravitons at (eb) the quantum critical point, they were able to identify and compute (e&eb) a set of perturbative amplitudes relevant for (e) black hole formation. These are the tree-level N-graviton scattering S-matrix elements in (eb) a kinematical regime (called classicalization limit) where (e) the two incoming ultra-Planckian gravitons produce (eb) a large number N of soft gravitons. Authors compute these amplitudes by using (e) the Kawai-Lewellen-Tye relations, as well as scattering equations and (e&eb) string theory techniques. Authors discover that this limit reveals (eb) the key features of the microscopic corpuscular black hole N-portrait. In particular, the perturbative suppression factor of a N-graviton final state, derived from (e) the amplitude, matches the non-perturbative black hole entropy when (e) N reaches the quantum criticality value, whereas final states with different value of N are (=) either suppressed or excluded by non-perturbative**

**corpuscular physics. Thus they identify the microscopic reason behind the black hole dominance over other final states including (e) non-black hole classical object. In the parameterization of the classicalization limit the scattering equations can be solved exactly allowing (eb) us to obtain closed expressions for the high-energy limit of the open and closed superstring tree-level scattering amplitudes for a generic number N of external legs. Authors demonstrate matching and complementarity between the string theory and (e&eb) field theory** in different large-s and large-N regimes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1016/j.nuclphysb.2015.02.004 Report number: MPP-2014-320, LMU-ASC 52/14 Cite as: arXiv:1409.7405 [hep-th] (or arXiv:1409.7405v2 [hep-th] for this version) **Black Hole Formation and Classicalization in Ultra-Planckian 2 -> N Scattering G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger STOP** Gravitational backgrounds, such as black holes, AdS, de Sitter and inflationary universes, should be viewed as (=) composite of N soft constituent gravitons. It then follows that such systems are close to (e) quantum criticality of graviton Bose-gas to Bose-liquid transition. Generic properties of the ordinary metric description, including (e) geodesic motion or particle-creation in the background metric, emerge as (=) the large-N limit of quantum scattering of (e) constituent longitudinal gravitons. **Gia Dvali, Cesar Gomez** show that this picture correctly accounts for (e) physics of large and small black holes in AdS, as well as reproduces (eb) well-known inflationary predictions for cosmological parameters. However, it anticipates new effects not captured by (e) the standard semi-classical treatment. In particular, **Gia Dvali, Cesar Gomez** predict observable corrections that are sensitive to (e&eb) the inflationary history way beyond last 60 e-foldings. They derive an absolute upper bound on the number of e-foldings, beyond (e) which neither de Sitter nor inflationary Universe can be approximated by (e) a semi-classical metric. However, they could in principle persist in a new type of (e) it quantum eternity state. Authors discuss implications of this phenomenon for (e&eb) the cosmological constant problem. Subjects: High Energy Physics - Theory (hep-th); Cosmology and

Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Quantum Physics (quant-ph) DOI: 10.1088/1475-7516/2014/01/023 Cite as: arXiv: 1312.4795 [hep-th] (or arXiv:1312.4795v1 [hep-th] for this version) **Quantum Compositeness of Gravity: Black Holes, AdS and Inflation Gia Dvali, Cesar Gomez STOP** As quantum information science approaches the goal of constructing quantum computers, understanding loss of information through (e&eb) decoherence becomes (=) increasingly important. The information about a system that can be obtained from (e) its environment can facilitate quantum control and error correction. Moreover, observers gain (eb) most of their information indirectly, by (e) monitoring (primarily photon) environments of (e) the "objects of interest." Exactly how this information is inscribed in the environment is essential for (e) the emergence of "the classical" from (e) the quantum substrate. In this paper, **Robin Blume-Kohout, Wojciech H. Zurek** examine how many-qubit (or many-spin) environments can store (e) information about a single system. The information lost to (e) the environment can be stored redundantly, or it can be encoded in (eb) entangled modes of the environment. Authors go on to show that randomly chosen states of the environment almost always encode (e) the information so that (e) an observer must capture (e) a majority of the environment to deduce (eb) the system's state. Conversely, in the states produced by (e) a typical decoherence process, information about a particular observable of the system is stored redundantly. This selective proliferation of "the fittest information" (known as Quantum Darwinism) plays a key role in (eb) choosing the preferred, effectively classical observables of (e) macroscopic systems. The developing appreciation that the environment functions not just as (=) a garbage dump, but as a communication channel, is extending understanding of the environment's role in (e&eb) the quantum-classical transition beyond (e) the traditional paradigm of decoherence. Keywords Decoherence quantum Darwinism quantum information quantum-classical transition **Foundations of Physics November 2005, Volume 35, Issue 11, pp 1857-1876 A Simple Example of "Quantum Darwinism": Redundant Information Storage in Many-Spin Environments Robin Blume-Kohout, Wojciech H. Zurek STOP Harold Ollivier, David**

**Poulin, and Wojciech H. Zurek** study the emergence of objective properties in open quantum systems. In our analysis, the environment is promoted from a passive role of a reservoir selectively destroying quantum coherence to an active role of amplifier selectively **proliferating information about the system. We show that only preferred pointer states of the system can leave a redundant and therefore easily detectable imprint on the environment. Observers who—as is almost always the case—discover the state of the system indirectly (by probing a fraction of its environment) will find out only about the corresponding pointer observable. Many observers can act in this fashion independently and without perturbing the system. They will agree about its state. In this operational sense, preferred pointer states exist objectively. Received 24 July 2003 DOI:http://dx.doi.org/10.1103/PhysRevLett.93.220401 Objective Properties from Subjective Quantum States: Environment as a Witness Harold Ollivier, David Poulin, and Wojciech H. Zurek Phys. Rev. Lett. 93, 220401 – Published 22 November 2004 STOP** Quantum Darwinism recognizes that we—the observers—acquire our information about the "systems of interest" indirectly from their imprints on the environment. **Michael Zwolak, H. T. Quan, and Wojciech H. Zurek** show that information about a system can be acquired from a mixed-state, or hazy, environment, but the storage capacity of an environment fragment is suppressed by its initial entropy. In the case of good decoherence, the mutual information between the system and the fragment is given solely by the fragment's entropy increase. For fairly mixed environments, this means a reduction by a factor 1−h, where h is the haziness of the environment, i.e., the initial entropy of an environment qubit. Thus, even such hazy environments eventually reveal the state of the system, although now the intercepted environment fragment must be larger by ~ (1−h)−1 to gain the same information about the system. DOI:http://dx.doi.org/10.1103/PhysRevLett.103.110402 **Quantum Darwinism in a Mixed Environment Michael Zwolak, H. T. Quan, and Wojciech H. Zurek Phys. Rev. Lett. 103, 110402 – Published 8 September 2009** We study quantum Darwinism—the redundant recording of information about the preferred states of a decohering system by its environment—for an object illuminated by a

blackbody. In the cases of point-source and isotropic illumination, we calculate the quantum mutual information between the object and its photon environment. We demonstrate that this realistic model exhibits fast and extensive proliferation of information about the object into the environment and results in redundancies orders of magnitude larger than the exactly soluble models considered to date. DOI:http://dx.doi.org/10.1103/PhysRevLett.105.020404 **Quantum Darwinism in an Everyday Environment: Huge Redundancy in Scattered Photons C. Jess Riedel and Wojciech H. Zurek Phys. Rev. Lett. 105, 020404 – Published 9 July 2010 STOP** Measurements transfer information about a system to the apparatus and then, further on, to observers and (often inadvertently) to the environment. I show that even imperfect copying essential in such situations restricts possible unperturbed outcomes to an orthogonal subset of all possible states of the system, thus breaking the unitary symmetry of its Hilbert space implied by the quantum superposition principle. Preferred outcome states emerge as a result. They provide a framework for "wave-packet collapse," designating terminal points of quantum jumps and defining the measured observable by specifying its eigenstates. In quantum Darwinism, they are the progenitors of multiple copies spread throughout the environment—the fittest quantum states that not only survive decoherence, but subvert the environment into carrying information about them—into becoming a witness. DOI:http://dx.doi.org/10.1103/PhysRevA.76.052110 **Quantum origin of quantum jumps: Breaking of unitary symmetry induced by information transfer in the transition from quantum to classical Wojciech Hubert Zurek Phys. Rev. A 76, 052110 – Published 16 November 2007 STOP** The Born rule, a foundational axiom used to deduce probabilities of events from **wavefunctions, is indispensable in the everyday practice of quantum physics. It is also key in the quest to reconcile the ostensibly inconsistent laws of the quantum and classical realms, as it confers physical significance to reduced density matrices, the essential tools of decoherence theory. Following Bohr's Copenhagen interpretation, textbooks postulate the Born rule outright. However, recent attempts to derive it from (e) other quantum principles have been successful, holding promise for (e)**

**simplifying and clarifying the quantum foundational bedrock. A major family of derivations is based on envariance, a recently discovered symmetry of (e) entangled quantum states. Here, authors identify and experimentally test three premises central to (e) these envariance-based derivations, thus demonstrating, (eb) in the microworld, the symmetries from (e) which the Born rule is derived. Further, they demonstrate envariance in a purely local quantum system, showing (eb) its independence from relativistic causality.Subjects: Quantum Physics (quant-ph)Cite as:arXiv:1604.01471 [quant-ph] (or arXiv:1604.01471v1 [quant-ph] for this version) Quantum probabilities from quantum entanglement: Experimentally unpacking the Born rule Jérémie Harris, Frédéric Bouchard, Enrico Santamato, Wojciech H. Zurek, Robert W. Boyd, Ebrahim Karimi**

Integrating $[m_4]_4(t, a_4)$ with respect to a from 0 to ∞, we have

$$\frac{dM_4(t)}{dt} = \int_0^\infty [mt]_4(t, a_4)da_4$$

$$= -\int_0^\infty [m_a]_4\,(t, a_4)da_4 - \int_0^\infty \mu(a_4)m(t, a_4)da_4$$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f^{'}{}_{4*}]_4 \int_0^\infty p_4^*(a_4)da_4\right)M_4 - [f_{2*}^4]_4 \int_0^\infty p_4^*(a_4)da_4 W_4$$

$= m_4(t, 0) - \int_0^\infty \mu_4(a_4)\, m_4(t, a_4)da_4$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}^{i}]_4 P_4^*\right)M_4 - [f_{2*}^{i}]_4 P_4^* W_{\cdot 4}$$

From (30), we arrive at the equations

$$\frac{dW_4(t)}{dt} = \int_0^{\infty} ([h_*]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4$$

$$-\left([f_{4*}] + [f_{2*}]_4 + \left[f_{4*}^{'}\right]_4 P_4^*\right) M_4 + \left(\frac{\left[h_*^{'}\right]_4}{[h_*]^4}\, P_4^*(0) - f_{2*}^{'} P_4^*\right) W_4,$$

(34)

$$\frac{dW_4(t)}{dt}\left[g_{*}^{'}\right]_4 Q_4^*\, M_{.4}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{r*}]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4 = ([h_*]_4\beta_4(\theta_4) - \mu_4(\theta_4))M_{,4}$$

Where $0< \theta_4 = \theta_4(t) < \infty$.

And if $a_4^2\,[a_4^2]_4<4\,[a_2]_4<1\,a_2 e$,

$\mathrm{Re}\,\lambda_4\left[\lambda_4\right]_4 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_4 = -\frac{4}{2}a_4[a2\ ]_4 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_*\,[h_*]_4\beta_4(a_4) - \mu_4(a_4)\right|$. Then$|n_4(t)|$ ⊠⊿∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_4^2 <4\, a_2 e$ and $[a_4] > 2\Delta$,or $[a_4^2]$ ⊠⊿$[a_2]e$ and $(a_4 - \sqrt{(a_4^2)_4 - 4(a_2)_4 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+4}\right)$ and $\beta_4(a_4) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4) + d\tau_4 da} = \frac{4}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_4 + [\mu_0]$ 0.5. Assume$[f_4](x_4)= K_4 x_4$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_4(x_4)= [0x_4]/(1 + x_4)$; the functional responses $[[f_2](x) = [0]\, x_4/(1 + x_4)$and $h_4(x_4)=(1+\omega_4 x_4)/(1 + x_4)$(See Holling [7] and Rashevsky [13])

Equation (44),now has the form of

$$\frac{m_4 + [\mu_0] + [K][P_0])}{[m_4]_4 + (\mu_0)_4 + [K][P_0])\frac{3}{4}\,[K]_4[P_0]\frac{3}{4}} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(4 + P_0) < b_4$ i e $[P_0] < b/(\delta_4 - b_4)$then the equilibrium$(P_0)$ $[P_0]_4(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_4$

$[f_{4*}] = K_4\, P^*\, [f_{2*}] = \frac{\sigma Q_4^*}{4+Q_4^*} \quad [h_*] = \frac{4+\omega_4 Q_4^*}{4+Q_4^*}$

$[f_{4*}] = K_4\, [f_{2*}^4] = \frac{\sigma_4}{(4+Q_4^*)^2}\, [g_*^4] = \frac{\sigma_4}{(4+P_4^*)^2} \quad [h_*^4] = \frac{\omega_4 - 4}{(4+Q_4^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0$, K, b, $\delta$, and $\omega$ need to satisfy

$$\frac{\omega_4}{(m + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma)^2} + \frac{2\omega_4}{([m_4]_4 + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma_4)^3} < 1$$

And

$$\frac{1}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^2}+\frac{2\omega_4}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4}+(m_4+[\mu_0]_4+\frac{K_4b_4}{\delta_4-b_4}+\sigma_4)\,Q_4^*)^3$$

$$=m_4+[\mu_0]_4+3+\frac{K_4b_4}{\delta_4-b_4}(4+Q^*)(1+\omega_4Q_4^*$$

Now

$$\left[a_4^2\right]-1\left[a_2\right]e=\left(\frac{2K_4b_4}{\delta_4-b_4}\right)^2+\frac{1K_4b_4\sigma_4\,Q_4^*}{(\delta_4-b_4)(1+Q_4^*)}+\frac{\sigma_4^2Q_4^{*2}}{(1+Q_4^*)^2}$$

$$\frac{1\,b_4\sigma_4\,(\delta_4-b)Q^*}{\delta_4(1+Q_4^*)^2}-\frac{1\,(\delta_4-b_4)^2(1-\omega_4)Q_4^*}{\delta_4(1+\omega_4Q_4^*)(1+Q_4^*)}P^*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_4\,ae_4^{mu}\;[\mu_0]_4-\frac{a_4}{a_4+4}$$

Hence if $[a_4^2>4(a_2)_4e$ and

$$[a_4] = \frac{2K_4 b_4}{\delta_4 - b_4} + \frac{\sigma_4 \, Q_4^*}{1 + Q_4^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_4]_4 = [f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}] P_4^*$$

$$[a_2]_4 = \left[f_{2*}^4\right] P_4^* - \frac{h_*^4}{h_*} P^*(0)$$

$$e = [g_*^4]_4 \, Q_4^*$$

$$\eta_4(t) = \left[h_*\right]_4 \beta_4 \left(\theta_4(t)\right) - \mu_4\left(\theta(t)\right),$$

The system(34)can be written as the following system:

$$[M(t)]\left[a_4\right] M_4 - [a_2]_4 W_4 - [n(t)] M_4$$

$$w_4(t) = eM_4$$

Write

$$A = \begin{pmatrix} -[a_4]_4 - [a_2]_4 & \\ e & 0 \end{pmatrix}$$

$$B=(t)= \begin{pmatrix} [\eta(t)]_4 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_4 \\ W_4 \end{pmatrix}$$

The solution of (32)is

$$X_4(t) = e^{A,t}\,[X]_4(0) + \int_0 e^{At}$$

**Some of the other conceptual and thematic and discursive factorial representationalities to which the model is applicable are:every one-dimensional quantum mechanical Hamiltonian H1 can have (e) a partner H2 such that H1 and H2 taken together may be viewed as (=) the components of a supersymmetric Hamiltonian. The term 'supersymmetric Hamiltonian' is taken to mean (eb) a Hamiltonian defined in terms of (e&eb)charges that obey (eb)the same algebra as that of (e) the generators of supersymmetry in field theory. The consequences of this symmetry for (er) the spectra of H1 and H2 are explored. It is shown how the supersymmetric pairing may be utilised to eliminate (e) the ground state of H1, or add a state below the ground state of H1 or maintain the spectrum of H1. It is also explicitly demonstrated that the supersymmetric pairing may be used to generate (eb) a class of anharmonic potentials with (e&eb) exactly specified spectra. The complete spectrum of an anharmonic potential so generated consists of (e) all the eigenstates of the simple harmonic oscillator and, in addition, a ground state at a specified energy E which lies arbitrarily below (eb) the E=1/2 ground state of the harmonic oscillator. Supersymmetric quantum mechanics of one-dimensional systems C V Sukumar Journal of Physics A: Mathematical and General, Volume 18, Number 15 STOP Authors examine the properties and consequences of (e&eb) pseudo-supersymmetry for (e) quantum systems admitting (e) a pseudo-Hermitian Hamiltonian. Authors explore the Witten index of (e) pseudo-**

**supersymmetry and show that every pair of diagonalizable (not necessarily Hermitian) Hamiltonians with (e&eb) discrete spectra and (e&eb) real or complex-conjugate pairs of (e) eigenvalues are (=) isospectral and have (e) identical degeneracy structure except (e) perhaps for the zero eigenvalue if and only if (e) they are pseudo-supersymmetric partners. This implies (eb) that pseudo-supersymmetry is (=) the basic framework for (e) generating non-Hermitian PT-symmetric and (e&eb) non-PT-symmetric Hamiltonians with (e&eb) a real spectrum via (e&eb) a Darboux transformation, and shows (eb) that every diagonalizable Hamiltonian H with(e&eb) a discrete spectrum and (e&eb) real or complex-conjugate pairs of (e) eigenvalues may be factored as (=) H=L♯L where L is (=) a linear operator with pseudo-adjoint L♯. In particular, this factorization applies to (e&eb) PT-symmetric and Hermitian Hamiltonians. The non-degenerate two-level systems provide (eb) a class of Hamiltonians that are pseudo-Hermitian. Authors demonstrate the implications of general results for (e&eb) this class in some detail. Copyright © 2002 Elsevier Science B.V. All rights reserved. Nuclear Physics B Volume 640, Issue 3, 30 September 2002, Pages 419–434Pseudo-supersymmetric quantum mechanics and isospectral pseudo-Hermitian Hamiltonians Ali Mostafazadeh STOP Michele Caselle, Alessandro Nada, Marco Panero present a high-precision lattice calculation of (e) the equation of state in the confining phase of (e) SU(2) Yang-Mills theory. Authors show that the results are described very well by (e) a gas of massive, non-interacting glueballs, provided (e) one assumes an exponentially growing Hagedorn spectrum. The latter can be derived within (eb) an effective bosonic closed-string model, leading to (eb) a parameter-free theoretical prediction, which is in perfect agreement with (=) proposed lattice results. Furthermore, when applied to (e&eb) SU (3) Yang-Mills theory, this effective model accurately describes (eb) the lattice results reported by Bors\'anyi et al. in JHEP 07 (2012) 056. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 07 (2015) 143 DOI: 10.1007/JHEP07 (2015)143 Cite as: arXiv:**

**1505.01106 [hep-lat] (or arXiv: 1505.01106v2 [hep-lat] for this version) Hagedorn spectrum and thermodynamics of SU (2) and SU (3) Yang-Mills theories Michele Caselle, Alessandro Nada, Marco Panero STOP B. de Wit J. Hoppe, H. Nicolai study the quantum-mechanical properties of (e) a supermembrane and examine the nature of (e) its ground state. A supersymmetric gauge theory of (e) area-preserving transformations provides (eb) a convenient framework for this study. The supermembrane can be viewed as (=) a limiting case of a class of models in supersymmetric quantum mechanics. Its mass does not depend on (=) the zero modes and vanishes only if (e) the wave function is (=) a singlet under (e&eb) supersymmetry transformations of (e) the nonzero modes. Authors exhibit (eb) the complexity of (e) the supermembrane ground state and (e&eb) examine various truncations of these models. None of these truncations has massless states. Copyright © 1988 Published by Elsevier B.V. Nuclear Physics B Volume 305, Issue 4, 28 December 1988: Pages 545-581: On the quantum mechanics of supermembranes Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. B. de Wit J. Hoppe, H. Nicolai STOP In this note we explain how world-volume geometries of (e) D-branes can be reconstructed within (eb) the microscopic framework where (e) D-branes are described through (e&eb) boundary conformal field theory. Authors extract the (non-commutative) world-volume algebras from (e) the operator product expansions of (e) open string vertex operators. For branes in a flat background with (e&eb) constant non-vanishing B-field, the operator products are computed perturbatively to (e) all orders in the field strength. Resulting series coincides with (=) Kontsevich's presentation of the Moyal product. After extending these considerations to (e) fermionic fields authors conclude with some remarks on the generalization of approach to (e&eb) curved backgrounds. D-branes and deformation quantization Volker Schomerus1,2 Published 25 July 1999 Journal of High Energy Physics, Volume 1999, JHEP06(1999)STOP Authors employ the path integral approach developed in Gangopadhyay S. and Scholtz F. E., Phys. Rev.**

**Lett., 102 (2009) 241602 to discuss the (generalized) harmonic oscillator in (eb) a noncommutative plane. The action for (e) this system is derived in (e&eb) the coherent state basis with (e&eb) additional degrees of freedom. From this the action in the coherent state basis without (e) any additional degrees of freedom is obtained. This gives (eb) the ground-state spectrum of the system. Authors then employ the exact renormalization group approach to show (eb) that equivalence can be constructed between this (noncommutative) system and (e&eb) a commutative system Path integral action and exact renormalization group dualities for quantum systems in noncommutative plane Sunandan Gangopadhyay Published 25 June 2015 • Copyright © EPLA, 2015 STOP Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini present a lattice study of (e)the equation of state in Yang-Mills theory based on (e) the exceptional G(2) gauge group. As is well-known, at zero temperature this theory shares (e&eb) many qualitative features with real-world QCD, including (e) the absence of colored states in the spectrum and dynamical string breaking at (e) large distances. In agreement with previous works, authors show that at finite temperature this theory features (eb) a first-order deconfining phase transition, whose nature can be studied by (e) a semi-classical computation. Authors also show that the equilibrium thermodynamic observables in the deconfined phase bear striking quantitative similarities with (=) those found in SU (N) gauge theories: in particular, these quantities exhibit (eb) nearly perfect proportionality to the number of gluon degrees of freedom, and (e&eb) the trace anomaly reveals (eb) a characteristic quadratic dependence on (e&eb) the temperature, also observed in SU (N) Yang-Mills theories (both in four and in three spacetime dimensions). Authors compare our lattice data with (e&eb) analytical predictions from (e) effective models, and discuss their implications for (e) the deconfinement mechanism and (e&eb) high-temperature properties of (e) strongly interacting, non-supersymmetric gauge theories. Results give (eb) strong evidence for (e) the conjecture that the thermal deconfining transition is governed by (e&eb) a universal mechanism, common to (e) all simple gauge groups.**

**Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 03 (2015) 057 DOI: 10.1007/JHEP03 (2015)057 Report number: DESY-14-146; IFT-UAM/CSIC-14-076 Cite as: arXiv: 1409.8305 [hep-lat] (or arXiv: 1409.8305v2 [hep-lat] for this version Exceptional thermodynamics: The equation of state of G(2) gauge theory Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini STOP In U (1) lattice gauge theory in three spacetime dimensions, the problem of confinement can be studied analytically in (e&eb) a semi-classical approach, in terms of (e&eb) a gas of monopoles with (e&eb) Coulomb-like interactions. In addition, this theory can be mapped to (e&eb) a spin model via (e&eb) an exact duality transformation, which allows (eb) one to perform high-precision numerical studies of (e) the confining potential. Taking advantage of these properties, authors carried out an accurate investigation of (e) the effective string describing (eb) the low-energy properties of flux tubes in (eb) this confining gauge theory. Authors found striking deviations from (e) the expected Nambu-Goto-like behavior, and, for the first time, evidence for (e) contributions that can be described by a term proportional to the extrinsic curvature of the effective string worldsheet. Such term is allowed by (e) Lorentz invariance and its presence in the infrared regime of the U (1) model was indeed predicted by (e) Polyakov several years ago. Results show that this term scales as expected according to (e) Polyakov's solution, and becomes (=) the dominant contribution to (e) the effective string action in the continuum limit. Authors also demonstrate analytically that the corrections to (e) the confining potential induced by (e) the extrinsic curvature term can be related to (e&eb) the partition function of (e) the massive perturbation of (e) a c=1 bosonic conformal field theory. The implications of results for SU (N) Yang-Mills theories in (e&eb) three and in four spacetime dimensions are discussed. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th) Journal reference: JHEP 1501 (2015) 105 DOI: 10.1007/JHEP01 (2015)105 Report number: IFT-UAM/CSIC-14-052 Cite as: arXiv:**

**1406.5127 [hep-lat] (or arXiv: 1406.5127v2 [hep-lat] for this version) A different kind of string Michele Caselle, Marco Panero, Roberto Pellegrini, Davide Vadacchino STOP In the past ten years, the ideas of supersymmetry have been profitably applied to (e&eb) many nonrelativistic quantum mechanical problems. In particular, there is now a much deeper understanding of why certain potentials are analytically solvable and (e&eb) an array of powerful new approximation methods for (e) handling potentials which are not (e) exactly solvable. In this report, authors review the theoretical formulation of (e) supersymmetric quantum mechanics and discuss many applications. Exactly solvable potentials can be understood in terms of (e&eb) a few basic ideas which include (e) supersymmetric partner potentials (e&eb) shape invariance and (e&eb) operator transformations. Familiar solvable potentials all have (e) the property of shape invariance. Authors describe new exactly solvable shape invariant potentials which include (e) the recently discovered self-similar potentials as (=) a special case. The connection between inverse scattering,(e&eb) isospectral potentials and (e&eb) supersymmetric quantum mechanics is discussed and multisoliton solutions of (e) the K dV equation are constructed. Approximation methods are also discussed within (eb) the framework of supersymmetric quantum mechanics and in particular it is shown that a supersymmetry inspired WKB approximation is (=) exact for a class of shape invariant potentials. Supersymmetry ideas give (eb) particularly nice results for the tunneling rate in (eb) a double well potential and for improving (eb+) large N expansions. Authors also discuss the problem of (e) a charged Dirac particle in (eb) an external magnetic field and other potentials in terms of (e&eb) supersymmetric quantum mechanics. Finally, authors discuss structures more general than supersymmetric quantum mechanics such as(=) parasupersymmetric quantum mechanics in which there is symmetry between a boson and (e&eb) a para-fermion of order p. Copyright © 1995 Published by Elsevier B.V. Physics Reports Volume 251, Issues 5–6, January 1995, Pages 267–385Supersymmetry and quantum mechanics Fred Coopera, Avinash Khareb, Uday Sukhatmec STOP The free energy of (e) the maximally**

**supersymmetric SU (N) gauge theory at temperature T is expected to scale, in the large-N limit, as (=) N2T4 times a function of the 't Hooft coupling, f (gYM2N). In the strong coupling limit the free energy has been deduced from (e) the near-extremal 3-brane geometry, and its normalization has turned out to be Full-size image (<1 K) times that found in (eb) the weak coupling limit. In this paper authors calculate the leading correction to (e) this result in inverse powers of (e) the coupling, which originates from (e) the R4 terms in the tree-level effective action of (e) type IIB string theory. The correction to Full-size image (<1 K) is positive and of (e) order Full-size image (<1 K). Thus, f (gYM2N) increases as (e&eb) the 't Hooft coupling is decreased, in accordance with (e&eb) the expectation that is should be approaching 1 in (eb) the weak coupling limit. Authors also discuss similar corrections for (e) other conformal theories describing (eb) coincident branes. In particular, authors suggest (eb) that the coupling independence of (e) the near extremal entropy for (e) D1-branes bound to (e) D5-branes is related to (e&eb) the vanishing of (e) the Weyl tensor of AdS3 × S3. Copyright © 1998 Published by Elsevier B.V. Nuclear Physics B Volume 534, Issues 1–2, 23 November 1998, Pages 202–222Coupling constant dependence in the thermodynamics of N = 4 supersymmetric Yang-Mills theory Steven S. Gubsera, Igor R. Klebanovb, 1, Arkady A. Tseytlinc, 2 STOP A new approach to (e) N=2 supersymmetry based on (e) the concept of harmonic superspace is proposed and is used to give (eb) an unconstrained superfield geometric description of N=2 super Yang-Mills and supergravity theories as (=) well as of matter N=2 hypermultiplets. The harmonic N=2 superspace has (e) an independent coordinates, in addition to the usual ones, the isospinor harmonics ui+or- on (e&eb) the sphere SU (2)/U (1). The role of ui+or- is to relate the SU (2) group realised on the component fields to (e&eb) a U (1) group acting on (e&eb) the relevant superfields. Their introduction makes it possible to SU (2)-covariantise (eb (t), e&eb) the notion of Grassmann analyticity. Crucial for the construction is the existence of (e) an analytic subspace of the general harmonic N=2 superspace. The hypermultiplet superfields and (e&eb) the true prepotentials (pre-**

**prepotentials) of N=2 super Yang-Mills and supergravity are (=) unconstrained superfunctions over (e&eb) this analytic subspace. The pre-prepotentials have (e) a clear geometric interpretation as gauge connections with (e&eb) respect to the internal SU (2)/U (1) directions. A radically new feature arises: the number of gauge and (e&eb) auxiliary degrees of freedom becomes (=) infinite while (e) the number of physical degrees of freedom remains (=) finite. Other new results are the massive N=2 Yang-Mills theory and (e&eb) various off-shell self-interactions of hypermultiplets. The propagators for matter and Yang-Mills superfields are given. Unconstrained N=2 matter, Yang-Mills and supergravity theories in harmonic superspace A Galperin, E Ivanov, S Kalitzin, V Ogievetsky and E Sokatchev STOP By assuming the existence of a novel multipronged string state for (e) D-particles interacting with (e&eb) D-brane intersections in type IIA string theory, authors are able to derive a quantum mechanical description of (e) supersymmetric Reissner-Nordstrom black holes. A supersymmetric index calculation provides (eb) evidence for (e) this conjecture. The quantum mechanical system becomes (=) two decoupled conformal quantum mechanical systems in (eb) the low energy limit. The conformal quantum mechanics has expected properties of (e) a dual description of string theory on (e&eb) AdS2×S2. Subjects: High Energy Physics - Theory (hep-th) Journal reference: JHEP 0001:006, 2000 DOI: 10.1088/1126-6708/2000/01/006 Report number: LBNL-43734 Cite as: arXiv:hep-th/9907101 (or arXiv:hep-th/9907101v5 for this version) Supersymmetric Quantum Mechanical Description of Four Dimensional Black Holes Julie D. BlumSTOP Using recently proposed soliton equations authors derive a basic identity for the scaling violation of (e) N=2 supersymmetric gauge theories Σiai∂F/∂ai−2F=8πib1u. Here F is (=) the prepotential, ai's are (=) the expectation values of (e) the scalar fields in (eb) the vector multiplet, u=1/2 Tr <ϕ2> and b1 is the coefficient of (e) the one-loop β-function. This equation holds in (eb) the Coulomb branch of all N=2 supersymmetric gauge theories coupled with (e&eb) massless matter. TOHRU EGUCHI and SUNG-KIL YANG, Mod Phys. Lett. A 11, 131 (1996).**

**DOI: 10.1142/S0217732396000151 PREPOTENTIALS OF N=2 SUPERSYMMETRIC GAUGE THEORIES AND SOLITON EQUATIONS STOP The Coulomb branch of N = 2 supersymmetric gauge theories in four dimensions is described in general by (e) an integrable Hamiltonian system in (eb) the holomorphic sense. A natural construction of such systems comes from (e) two-dimensional gauge theory and (e&eb) spectral curves. Starting from this point of view, authors propose an integrable system relevant to (e&eb) the N = 2 SU (n) gauge theory with (e&eb) a hypermultiplet in the adjoint representation, and offer (eb) much evidence that it is correct. The model has (e) a Full-size image (<1 K) S-duality group (with the central element −1 of Full-size image (<1 K) acting as (=) charge conjugation); Full-size image (<1 K) permutes the Higgs, confining, and oblique confining phases in the expected fashion. We also study more exotic phases. Copyright © 1996 Published by Elsevier B.V. Nuclear Physics B Volume 460, Issue 2, 5 February 1996, Pages 299–334Supersymmetric Yang-Mills theory and integrable systems Ron Donagia, 1, Edward Wittenb, 2 STOP A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 present a first step towards generalizing (eb) the work of Seiberg and Witten on N = 2 supersymmetric Yang-Mills theory to (e&eb) arbitrary gauge groups. Specifically, authors propose a particular sequence of (e) hyperelliptic genus n − 1 Riemann surfaces to underly (e&eb) the quantum moduli space of SU (n)N = 2 supersymmetric gauge theory. These curves have (e) an obvious generalization to (e) arbitrary simply laced gauge groups, which involves (e&eb) the A-D-E type simple singularities. To support proposal, authors argue that the monodromy in (e&eb) the semiclassical regime is (=) correctly reproduced. We also give some remarks on a possible relation to string theory. Copyright © 1995 Published by Elsevier B.V. Physics Letters B Volume 344, Issues 1–4, 26 January 1995, Pages 169–175Simple singularities and N = 2 supersymmetric Yang-Mills theory A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 STOP Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 review and extend earlier work that uses (e) the AdS/CFT correspondence to relate the**

**black-hole–black-string transition of gravitational theories on a circle to (e&eb) a phase transition in maximally supersymmetric (1 + 1)-dimensional SU(N) gauge theories at (eb) large N, again compactified on (e&eb) a circle. Authors perform gravity calculations to determine (eb) a likely phase diagram for (e) the strongly coupled gauge theory. Authors then directly study the phase structure of (e) the same gauge theory, now at (eb) weak't Hooft coupling. In the interesting temperature regime for (e) the phase transition, the (1 + 1)-dimensional theory reduces to (e&eb) a (0 + 1)-dimensional bosonic theory, which they solve using (e) Monte Carlo methods. Authors find strong evidence that the weakly coupled gauge theory also exhibits (eb) a black hole–black string-like phase transition in (eb) the large N limit. Authors demonstrate that a simple Landau–Ginzburg-like model describes (eb) the behaviour near the phase transition remarkably well. The weak coupling transition appears to be (=) close to the cusp between a first-order and a second-order transition. Black-hole–black-string phase transitions in thermal (1 + 1)-dimensional supersymmetric Yang–Mills theory on a circle Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 Published 29 October 2004 • 2004 IOP Publishing Ltd Classical and Quantum Gravity, Volume 21, Number 22STOP Monitoring the orbits of stars around (e&eb) Sgr A* offers(eb)the possibility of detecting (e&eb) the precession of (e) their orbital planes due to (e) frame dragging, of (e) measuring the spin and quadrupole moment of (e) the black hole, and of (e)testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and (e&eb) tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are (=), in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for which an optical interferometer, such as GRAVITY, will detect orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to (e&eb) the timescale of precession due to (e) the quadrupole moment of (e) the black hole. As a result, this non-conservative effect is (=) a**

**potential source of systematic uncertainty in (eb) testing the no-hair theorem with (e&eb) stellar orbits. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI:10.1088/0004-637X/777/1/57 Cite as:arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this version) Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP Authors perform a phase space analysis of (e) a generalized modified gravity theory with (e&eb) non-minimal coupling between geometry and (e&eb) matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that (eb) in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, Authors prove that there is (=) an attractor solution to describe (eb) the late time accelerating universe when (e) the modified gravity is chosen in (eb) a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and supernova type Ia data sets, Authors find that (eb) the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as:arXiv: 1512.09281 [gr-qc] (or arXiv: 1512.09281v1 [gr-qc] for this version) Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang study the P−V criticality and phase transition in (eb) the extended phase space of (e) charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of the system and (e&eb) its conjugate quantity is (=) the thermodynamic volume of (e) the black hole. The black**

**holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb)no P−V criticality and phase transition appear, while for the black holes with a spherical horizon, even when (e) the charge of the black hole is absent, the P−V criticality and the small black hole/large black hole phase transition (eb) will appear, but it happens only in (eb) d=5 dimensions; when (e) the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to (e) have the P−V criticality and the small black hole/large black hole phase transition, there exists (eb) an upper bound for (e) the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the Gauss-Bonnet coefficient and Q is the charge of the black hole. Authors calculate the critical exponents at the critical point and find that for all cases, they are (=) the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09 (2013)005 Cite as: arXiv: 1306.6233 [gr-qc] (or arXiv: 1306.6233v4 [gr-qc] for this version) P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP paper studies the whole process of (e) gravitational collapse and (e&eb) accretion of (e) a massless scalar field in (eb) asymptotically flat spacetime. Two kinds of initial configurations are considered. One is the initial data without (e) black hole; the other contains (e) a black hole. Under suitable initial conditions, Authors find that multi-horizon will appear which mean (eb) that the initial black hole formed by (e) gravitational collapse or existing at the beginning will instantly expand and suddenly grow (e&eb, eb), rather than grows gradually in the accretion process. A new type of critical behavior is found around (e&eb) the instant expansion. The numerical computation shows (eb) that the critical exponents are (=) universal. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Computational Physics (physics. Comp-ph) Cite as: arXiv: 1604.03363 [gr-qc] (or arXiv: 1604.03363v1 [gr-qc] for this version)**

**Multi-horizon and Critical Behavior in Gravitational Collapse of Massless Scalar Zhoujian Cao, Rong-Gen CAI, Run-Qiu Yang STOP In this paper, Authors present the results on the gravitational collapse of (e) charged massless scalar field in asymptotically flat spacetime with(e&eb) a perfectly reflecting wall. Differing from previous works, Authors study the system in (eb) the double null coordinates, by which we could simulate (e&eb) the system until the black hole forms with (e&eb) higher precision but less performance time. Authors investigate the influence of charge on (e&eb) the black hole formation and the scaling behavior near the critical collapses. The gapless and gapped critical behaviors for (e) black hole mass and charge are studied numerically. Authors find that they satisfy (eb) the scaling laws for critical gravitational collapse but the gapped critical behavior is different from (e) its AdS counterpart. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1602.00112 [gr-qc] (or arXiv: 1602.00112v1 [gr-qc] for this version) Multiple critical gravitational collapse of charged scalar with reflecting wall Rong-Gen CAI, Run-Qiu Yang STOP In this paper Authors investigate the (2+1)-dimensional topological non-Fermi liquid in (eb) strongly correlated electron system, which has (e) a holographic dual description by (e) Einstein gravity in (3+1)-dimensional anti-de Sitter (AdS) space-time. In a dyonic Reissner-Nordstrom black hole background, Authors consider a Dirac fermion coupled to (e&eb) the background U (1) gauge theory and an (e&eb) intrinsic chiral gauge field bM induced by (e) chiral anomaly. UV retarded Green's function of (e) the charged fermion in (eb) the UV boundary from (e) AdS4 gravity is calculated, by imposing (e&eb) in-falling wave condition at the horizon. Authors also obtain IR correlation function of (e) the charged fermion at (eb) the IR boundary arising from (e) the near horizon geometry of (e) the topological black hole with (e&eb) index k=0,±1. By using (e) the UV retarded Green's function and IR correlation function Authors analyze the low frequency behavior of (e) the topological non-Fermi liquid at zero and (e&eb) finite temperatures, especially the relevant non-Fermi liquid behavior near (eb) the quantum critical region. In addition, Authors**

**find that when k=±1, the effective mass of (e) bulk Dirac fermion at IR fixed point is (=) topologically quantized, consequently the Fermi momentum presents (eb) a discrete sparse distributed pattern, in contrary to (e&eb) the Ricci flat case with k=0. As demonstrated examples, Authors calculate the spectral functions of (e) topological non-Fermi liquid due to (e) anomalous Hall effect (AHE) (|b→|≠0,b 0=0) and chiral magnetic effect (CME) (|b→|=0,b0≠0), which can be tested in topological materials such as (=) topological insulator and (e&eb) topological metal with (e&eb) strong couplings. Subjects: High Energy Physics - Theory (hep-th); Strongly Correlated Electrons (cond-mat.str-el); General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv: 1601.03865 [hep-th] (or arXiv: 1601.03865v1 [hep-th] for this version) Topological Non-Fermi Liquid Rong-Gen Cai, Yong-Hui Qi, Yue-Liang Wu, Yun-Long Zhang STOP The unusual magnetic materials are significant in both science and technology. However, because of (e) the strongly correlated effects, it is difficult to understand their novel properties from (e) theoretical aspects. Holographic duality offers (eb) a new approach to understanding such systems from (e) gravity side. This paper will give a brief review of Authors recent works on the applications of holographic duality in understanding (e&eb) unusual magnetic materials. Some quantitative compare between holographic results and (e&eb) experimental data will be shown and some predictions from (e) holographic duality models will be discussed. Comments: For Second LeCosPA Symposium at Taiwan National University, Taipai, Dec. 14-18, 2015 Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1601.02936 [hep-th] (or arXiv:1601.02936v1 [hep-th] for this version) Understanding strongly coupling magnetism from holographic duality Rong-Gen Cai, Run-Qiu Yang STOP It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a**

**Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3**

**[astro-ph.CO] for this version) Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang**

# Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

## Module Number Six

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-465) 1990. We demonstrate its applicability cosmology/ quantum computation economics etal. For the first time, physicists couple distant nuclear spins using a single electron July 11, 2016. Some sentences have been deleted due to spatial restrictions. Please excuse me. For the first time, researchers at the University of Basel in Switzerland have coupled the nuclear spins of distant atoms using (e) just a single electron. Three research groups took part in this complex experiment, the results of which have now been published in the journal Nature Nanotechnology. In most materials, the nuclear spins of (e) neighboring atoms have (e) only a very weak effect on one another, as (=) the tiny nuclei are located deep within (e&eb) the atoms. This is not the case in metals, however, where (e) some of the electrons are able to move freely. The electron spins are able to couple nuclear spins at (eb) relatively large distances from one another. Named after four physicists, this RKKY interaction was discovered in the 1950s. Nuclear spins linked by (e&eb) an individual electron spin Now, an experiment by researchers from the Department of Physics at the University of Basel has for the first time succeeded in demonstrating this mechanism with (e&eb) a single electron, describing it using (e) quantum theory. The team led by Prof. Richard Warburton trapped a single electron inside (eb) a quantum dot. With use of a method developed in Basel to measure (eb) nuclear spin resonance, they showed that the electron coupled (e&eb) the spins of nuclei up to five nanometers apart - a huge distance in the world of (e) nuclear spins. The results are particularly relevant to (e&eb) the development of spin qubits: these seek to use (e) electron spins to carry (e&eb) information, but the interaction with (e&eb) the nuclei limits (e) the stability of the quantum information. Concentrated blast of physics More information: Role of the electron spin in determining the coherence of the nuclear spins in a quantum dot, Nature Nanotechnology, DOI: 10.1038/nnano.2016.114 Journal reference: Nature Nanotechnology search and more info website Provided by: University of Basel Read more at: http://phys.org/news/2016-07-physicists-couple-distant-nuclear-electron.html#jCp A new spin on quantum computing: Scientists train electrons with microwaves February 15, 2016 A silicon sample doped with (e&eb) bismuth atoms (left image) that is (=) just 150 nanometers thick is fitted with (e&eb) a superconducting resonator that includes (e) a capacitor (black, in left image; light gray in center image) and an inductive wire (red line … In what may provide (eb) a potential path to processing information in (eb) a quantum computer, researchers have switched an intrinsic property of electrons from an excited state to (e&eb) a relaxed state on demand using (e) a device that served as (=) a microwave "tuning fork." The team's findings could also lead to (eb) enhancements in magnetic resonance techniques, which are (=) widely used to explore (e) the structure of materials and biomolecules, and for medical imaging. The**

**international research team, which included scientists at the Department of Energy's Lawrence Berkeley National Laboratory (Berkeley Lab), demonstrated how to dramatically increase the coupling of (e) microwaves in a specially designed (e&eb) superconducting cavity to (e) a fundamental electron property called spin—which, like a coin, can be flipped. By zapping an exotic silicon material developed at Berkeley Lab with the microwaves, they found that they could rapidly change (e&eb) the electron spins from (e) an excited state to a relaxed, ground state by (e) causing (eb) the electrons to emit (eb) some of their energy in the form of (e) microwave particles known as photons. Left on their own, the electron spins would be extremely unlikely to flip back to (e) a relaxed state and to also emit (eb) a photon - the natural rate for this light-emitting effect, known as the Purcell effect, is about once every 10,000 years. The experiment demonstrated an accelerated, controllable relaxation of (e) electron spins and the release of (e) a microwave photon in about 1 second, said Thomas Schenkel, a physicist in Berkeley Lab's Accelerator Technology and Applied Physics Division who led the design and development of the silicon-bismuth sample used in the experiment. A new spin on quantum computing: Scientists train electrons with microwaves Like the classical tuning fork slung from guitar strings in this photo, scientists used (e) a device that worked like a microwave tuning fork to switch (e&eb) a fundamental property of electrons on (e&eb) demand in a silicon-bismuth sample developed at Berkeley Lab. Credit: Flickr/Pierre Guinoiseau "It's like a juggler who throws the balls up, and the balls come down 1,000 times faster than normal, and they also emit (eb) a microwave flash as they drop," he said. The results were published online Feb. 15 in the journal Nature. "Our results are highly significant for (e&eb) quantum information processing," said Patrice Bertet, a quantum electronics scientist at the French Atomic Energy Commission (CEA) who led the experiment. "Indeed, they are a first step toward (e&eb) the strong coupling of (e) individual electron spins to microwave photons, which could form (eb) the basis of new spin-based quantum computer architecture." John Morton, a professor at the London Center for Nanotechnology and co-author of the study, said, "Our ultimate aim is to find a link between quantum information that is fixed and (e&eb) quantum information that can be transported by photons." In today's computers, information is stored as (=) individual bits, and each bit can either be (=) a one or a zero. Quantum computers, though, could conceivably be (=) exponentially more powerful than modern computers because (e) they would use a different kind of bit, called a qubit, which because of (e) the weird ways of quantum mechanics can simultaneously behave (e&eb) as both a one and a zero. A coupled array of qubits would allow (eb) a quantum computer to perform (e&eb) many, many calculations at (eb) the same time, and electron spins are (=) candidates for qubits in (eb) a quantum computer. The latest study shows how the microwave photons could work in concert with (e&eb) the spins of electrons to move information in a new type of computer. "What we need now is ways to wire up these systems—to couple (e&eb) these spins together," Morton said. "We need to make coupled qubits that can perform (e&eb) computations." In the experiment, conducted at CEA in France, a small sample of a highly purified form of (e) silicon was implanted with (e&eb) a matrix of bismuth atoms, and a superconducting aluminum circuit was deposited on (e&eb) top to create (eb) a high-quality resonant cavity that allowed (eb) precise tuning of the microwaves. The electron spins of the bismuth atoms were then flipped into (e&eb) the excited, "spin-up" state. The microwave cavity was then tuned (e&eb) like a musical tuning fork, to (e) a particular resonance that coaxed the spins into emitting (eb) a photon as they flipped back to (e) a relaxed state. The cavity boosted the number of states into (e&eb) which a photon can be (=) emitted, which greatly increased (eb) the decay rate for (e) the electron spins in (eb) a controllable way. The technique is much like buying more lottery tickets to increase (eb) your chances of winning, Morton said. A new spin on quantum computing: Scientists train electrons with microwaves From left are, Berkeley Lab scientists Thomas Schenkel, Qing Ji and Peter Seidl at the NDCX-II (Neutralized Drift Compression Experiment II), which produces (eb) powerful ion beams. Researchers are exploring how to use NDCX-II to process exotic … The large bismuth atoms embedded in (eb) the silicon sample provided (eb) the electrons with (e&eb) unique spin properties that enabled (eb) the experiment. Schenkel said that implanting the bismuth atoms into (e&eb) the delicate silicon framework, a**

**process known as "doping," was "like squeezing bowling balls into (e&eb) a lattice of ping-pong balls." "We did a new trick with silicon. People wouldn't expect you could squeeze (e) anything new out of (e) silicon," Schenkel said. "Now we're looking into further improving (eb) bismuth-doped silicon and into (e&eb) tailoring the spin properties of (e) other materials, and using (e) this experimental technique for these materials." To enhance the performance of materials used (e) in future experiments, Schenkel said it will be necessary to improve (eb) the doping process so it is less damaging to (e) the silicon lattice. Also, the implantation process could be designed to produce (eb) regularly spaced arrays of (e) individual electron spins that would be more useful for (e) quantum computing than a concentrated ensemble of (e) electron spins. "We are now doing experiments on processing (e&eb) this and other materials at higher temperature and pressure with nanosecond ion pulses at NDCX-II, one of the accelerators here at Berkeley Lab," Schenkel said. "There are indications that it will improve (eb) the overall spin quality." Researchers said the latest research could potentially prove useful in boosting (eb) the sensitivity of scientific techniques like nuclear magnetic resonance spectroscopy and dynamic nuclear polarization, useful for a range of experiments, and could also shorten (e) experimental times by manipulating (e&eb) spin properties. "You need a way to reset (e&eb) spins—the ability to cause (eb) them to relax on demand to improve (eb) the rate at which you can repeat an experiment," Morton said. Bertet said it may be possible to further accelerate (eb) the electron-flipping behavior to (e) below 1 millisecond, compared to (e) the 1-second rate in the latest results. "This will then open the way to many new applications," he said. Explore further: Bismuth provides perfect dance partners for quantum computing qubits More information: A. Bienfait et al. Controlling spin relaxation with a cavity, Nature (2016). DOI: 10.1038/nature16944 Nature | Letter Controlling spin relaxation with a cavity A. Bienfait, J. J. Pla, Y. Kubo,X. Zhou, M. Stern, C. C. Lo, C. D. Weis, T. Schenkel, D. Vion,D. Esteve, J. J. L. Morton& P. Berte 531, 74–77 (03 March 2016) doi:10.1038/nature16944 Spontaneous emission of radiation is (=) one of the fundamental mechanisms by which an excited quantum system returns to (e) equilibrium. For spins, however, spontaneous emission is (=) generally negligible compared to (e&eb) other non-radiative relaxation processes because of (e) the weak coupling between the magnetic dipole and (e&eb) the electromagnetic field. In 1946, Purcell realized that the rate of spontaneous emission can be greatly enhanced by (e) placing the quantum system in a resonant cavity. This effect has since been used extensively to control (e&eb) the lifetime of atoms and semiconducting heterostructures coupled to (e&eb) microwave2 or optical3, 4 cavities, and is essential for (e) the realization of high-efficiency single-photon sources5. Here authors report the application of this idea to (e&eb) spins in solids. By coupling donor spins in silicon to (e&eb) a superconducting microwave cavity with (e&eb) a high quality factor and a small mode volume, authors reach the regime in which spontaneous emission constitutes (e) the dominant mechanism of spin relaxation. The relaxation rate is increased by (e) three orders of magnitude as the spins are tuned to (e) the cavity resonance, demonstrating that energy relaxation can be controlled (e&eb) on demand. Results provide a general way to initialize (eb) spin systems into their ground state and therefore have applications in (eb) magnetic resonance and (e&eb) quantum information processing6. They also demonstrate that the coupling between the magnetic dipole of a spin and (e&eb) the electromagnetic field can be enhanced up to (e) the point at which quantum fluctuations have (e) a marked effect on (e&eb) the spin dynamics; as such, they represent (eb) an important step towards the coherent magnetic coupling of (e) individual spins to (e) microwave photons Scientists obtain evidence of many-body localization in a closed quantum system July 8, 2016 by Olivia Meyer-Streng Science Cover Vol 352, Issue 6293, 24 June 2016, to highlight the research published in this issue (see original publication). The cover shows an artist's view of the experiment, illustrating the persistence of the density step of (e) an atomic distribution in an optical lattice. During equilibration ordinary many-body systems lose (e) all information about (e&eb) the initial state. Every morning authors experience an example for (e) this behaviour. Milk poured into a cup of coffee mixes (e&eb) perfectly and after some time it is impossible to say how exactly the two fluids were (=) put together. The same behaviour holds for (e) almost all quantum systems. However,**

**recently a new phenomenon called "many-body localization" has been predicted theoretically, which allows (eb) well insulated quantum systems to (e) preserve memory of the initial state forever. Now a team of scientists around Dr. Christian Groß and Professor Immanuel Bloch (Director at MPQ and Chair of Quantum Optics at LMU Munich), in cooperation with David Huse (Princeton University), has obtained evidence of such a behaviour in a two-dimensional quantum system of cold rubidium atoms trapped in (e&eb) an optical lattice. The scientists observed that – beyond a certain degree of disorder imprinted on (eb) the particle ensemble in the beginning – the system would relax into (e&eb) a steady state still containing (e) detailed microscopic information about its past. "We were able to observe the transition from a thermalized state into (e&eb) a many-body localized phase", Christian Groß points out. "It is the first observation of that kind in a regime that is not accessible with (e&eb) state-of-the-art simulations on classical computers." The experiment is not only of fundamental interest; the results might also lead to (eb) new ways for storing quantum information. Motivated by the foundational problem of how interacting particles behave (e&eb) in a disordered system, in the 1950s the American physicist Philip Warren Anderson discovered the famous localization phenomenon for (e) non-interacting particles, now called "Anderson localization". Here, disorder prevents (e) the particles to move and consequently all transport is (=) stopped. But what happens when disorder comes together with (e&eb) interactions? Will interactions lead to (eb) transport and thermalization, or will the localization persist even at (eb) high energies? So far, there is no theoretical model that faithfully predicts (eb) the evolution of a closed quantum system in more than one dimension under (e&eb) these conditions, although, the possibility for (e)localization has been theoretically suggested. In order to investigate these questions experimentally, strict requirements on controllability and perfect isolation of (e) the system have to be fulfilled. In the experiment described here, ultracold rubidium atoms are loaded into (eb) an optical lattice, a microscopic array of light traps formed by interfering (e&eb) laser beams. The disorder is created by projecting (e&eb) a computer-generated random light pattern onto (e&eb) the optical lattice. As a result of this "disorder", the depth of each microscopic trap in the lattice potential varies from site to (e) site. In fact, Prof. Bloch's group has advanced the experimental tools to such perfection that they can steer the position of (e) the atoms in the artificial light crystal and (e&eb) the interaction between (e&eb) them almost at will. With a high resolution microscope that detects the emitted fluorescence light, the position of each atom can be observed with (e&eb) high precision. Additionally, the atomic density distribution of the initial state can be (=) controlled, and for variable time intervals the evolution of the distribution can be (=) measured with high precision. These tools at hand, non-thermalizing behaviour can be probed in (eb) a conception ally very simple way. Any thermalized state of a closed system reflects (e&eb) the symmetry of its container – like, for example, water spilled into (e&eb) a round bowl would immediately cover (e) the whole bottom. In analogy, the scientists generate (eb) a density step in the initially prepared sample by (e) blowing away half of the atomic distribution with laser radiation. Then they watch how the remaining particles migrate into (e&eb) the empty half. For small disorders, the initially prepared density step is smeared out (e) fast and the initially empty and filled halves become (=) indistinguishable. However, when the measurement is repeated for (e) strong disorder, traces of the initial state remain and the system does not (e) relax to a thermal state even for (e) very long times. "We observe a fairly sharp onset of non-thermalizing behaviour above (e) a critical value", says Christian Groß. "This absence of thermalization is remarkable because (e) thermalization persists in a system of interacting particles, even at (eb) the high energies probed in the experiment." The scientists interpret their observation as (=) the onset of many-body localization in (eb) the atomic system. This is of fundamental interest because it means (eb) the breakdown of equilibrium statistical mechanics. On the other hand, the persistence of initial state information could be used as (=) a source for quantum information technologies. "It should also be emphasized that we obtain these results for a system size that is far beyond (e) numerically accessible scales", says Jae-yoon Choi, postdoc at the experiment. More information: J.-y. Choi et al. exploring the many-body localization transition in two dimensions, Science (2016). DOI: 10.1126/science.aaf8834 Journal reference: Science search and more**

**info website Provided by: Max Planck Society search and more info website Bosons refusing to thermalize in 2D Messy, interacting quantum-mechanical systems are difficult to analyze theoretically. In a single spatial dimension, the calculations are (=) still tractable, and experiments have recently confirmed (eb) the prediction that sufficiently strong disorder can disrupt (e&eb) the transport of interacting particles. In two dimensions, however, the theoretical blueprint is missing. Choi et al. used (e) single-site imaging of cold 87Rb atoms in (eb) an optical lattice to show (eb) that similar localization occurs in (eb) two-dimensional (2D) systems. The study highlights the power of quantum simulation to solve problems that are currently inaccessible to classical computing techniques. Science, this issue p. 1547 Abstract A fundamental assumption in statistical physics is that generic closed quantum many-body systems thermalize under (e&eb) their own dynamics. Recently, the emergence of many-body localized systems has questioned (e) this concept and challenged our understanding of the connection between statistical physics and (e&eb) quantum mechanics. Here authors report on the observation of a many-body localization transition between thermal and (e&eb) localized phases for (e) bosons in a two-dimensional disordered optical lattice. With our single-site–resolved measurements, authors track the relaxation dynamics of (e) an initially prepared out-of-equilibrium density pattern and find strong evidence for (e) a diverging length scale when (e) approaching the localization transition. Experiments represent a demonstration and in-depth characterization of many-body localization in (eb) a regime not accessible with (e&eb) state-of-the-art simulations on classical computers. Pelican instrument provides crucial experimental evidence of unusual quantum state July 8, 2016 Artist's illustration. Credit: Australian Nuclear Science and Technology Organisation (ANSTO) Experiments using inelastic neutron scattering at the Australian Centre for Neutron Scattering have found indications (eb) of a possible new quantum spin state in a novel antiferromagnetic material barium ytterbium zinc oxide (Ba3Yb2Zn5O11) which provides (eb) both a challenge and validation of (e) the third law of thermodynamics. The quantum mechanical information gained from (e) the experiment is crucial not only in understanding the quantum state of the specific material but also provides (eb) better understanding of high temperature superconductivity and ultimately offers (eb) guidance in the design of materials for (e) quantum computing. A research team led by Associate Professor Takatsugu Masuda at the Institute for Solid State Physics at the University of Tokyo with collaborators from ANSTO, Osaka University and Taiwan's National Cheng Kung University, have performed a successful experiment using Pelican, time-of-flight cold neutron spectrometer at ANSTO. The results have just been published in Physical Review B as a rapid Communication and an Editors' Suggestion. "The great success of this challenging experiment relied on (e) the excellent performance of the Pelican instrument, wonderful technical support at ANSTO and strong expertise in the relevant field from A/Prof Masuda's group," said Dr Dehong Yu and Dr Richard Mole, co-authors and Pelican instrument scientists. Pelican instrument provides crucial experimental evidence of unusual quantum state INS spectra measured at 1.5 K (a), 6 K (b), 12 K (c), and 40 K (d) using setup I. (e) INS spectrum measured at 1.5 K using the setup II. (f) Q dependence of the integrated intensity obtained from spectrum in (e). Red solid curve is the …….(deleted due to spatial constraints) Barium ytterbium zinc oxide (Ba3Yb2Zn5O11) is a novel antiferromagnetic model system having (e) a "breathing pyrochlore lattice" consisting of (e) arrays of alternating large and (e&eb) small corner-sharing tetrahedra of ytterbium atoms. Earlier research had suggested (eb) the material was (=) an isotropic spin half (S= ½) system. However, four low energy excitations measured on the sample with the Pelican instrument at a temperature of 1.5 K clearly disagree with (e) this designation and the associated theoretical model based on (e) the isotropic spins having simple interactions. "These observations suggested (eb) a much more complex spin state in the system," said Yu. "The experiment is interesting and important because as you start going very low in temperature, quantum mechanics takes over (e) system rather than thermodynamics, and you observe unusual properties," said Mole. "According to low temperature heat capacity measurements, the material does not (e) have zero entropy or a single ground state until you reach 63 milli Kelvin or -272.94 Celsius, before it starts obeying the third law." Pelican instrument provides crucial experimental evidence of (e) unusual quantum state**

**Breathing pyrochlore lattice formed by (e) A3+ ions as determined by Kimura, Nakatsuju and Kimura in previous work The third law of thermodynamics states that for pure crystals as the temperature of a system approaches absolute zero K, (eb) its entropy or disorder approaches zero. "The research clearly indicates (eb) that below this temperature the material may be in (eb) an exotic state, such as (=) a liquid-like spin state, in which the arrangement of spins behaves (e&eb) like a liquid. In this exotic state, novel phenomena may show up and be detectable; for example, charge and spin separation of electrons may be (=) possible," said Yu. The researchers suggest that further investigations, experimentally and theoretically, are required to ultimately determine the true identity of (e) the ground state of the material. "Discovering the existence of regular tetrahedral spin made it well worth coming all the way to Australia to make use of this neutron experimental apparatus," said Associate Professor Masuda. Many Japanese researchers are using ANSTO's neutron scattering instruments through the peer review system, while their nuclear facilities are offline. The outcome of this particular experiment is also part of the collaboration agreement between ANSTO and the University of Tokyo. More information: T. Haku et al. Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet, Physical Review B (2016). DOI: 10.1103/PhysRevB.93.220407 Journal reference: Physical Review B search and more info website: provided by: Australian Nuclear Science and Technology Organisation (ANSTO) search and more info website Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 T. Haku, K. Kimura, Y. Matsumoto, M. Soda, M. Sera, D. Yu, R. A. Mole, T. Takeuchi, S. Nakatsuji, Y. Kono, T. Sakakibara, L.-J. Chang, and T. Masuda Phys. Rev. B 93, 220407(R) – Published 24 June 2016 Authors study low-energy excitations in (eb) the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 by (e) a combination of inelastic neutron scattering (INS) and (e&eb) thermodynamical property measurements. The INS spectra are quantitatively explained by (e) spin-12 single-tetrahedron model having (e) XXZ anisotropy and Dzyaloshinskii-Moriya interactions. This model has a twofold degeneracy of (e) the lowest-energy state per tetrahedron and well reproduces (eb) the magnetization curve at (eb) 0.5 K and heat capacity above 1.5 K. At lower temperatures, however, authors observe a broad maximum in (eb) the heat capacity around 63 mK, demonstrating (eb) that a unique quantum ground state is selected due to (e) extra perturbations with (e&eb) an energy scale smaller than the instrumental resolution of INS. Revised 13 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.220407 © 2016 American Physical Society Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA (Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor. In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface. As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone. Such planets are all the "right" distance from (e) their star, and hopefully, also have water. But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone. For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement. But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway. But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain. It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of**

**reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life. Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface. The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113 Abstract The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question. Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry. This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry. Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments. Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy. The electronic and (e&eb) structural variability among them is extraordinary. The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations. The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface. The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers. Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp**

## 48. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_6}{dt} = -b_6 Q_6 + g_6(g_6)Q_6$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_6(0) = 0, \quad g_6(P_6) > 0, \quad P_6 \ 0, \qquad \lim_{P_6 \to \infty} g_6(P_6) = (K_6)_6 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_6 + (\rho_a)_6 + D_6(a_6, P_6, Q_6)\rho_6 + 0, \qquad 0 < a_6 \ \infty,$$

$$\rho_6(t, 0) = \int_0^\infty n_6(a_6, P_6, Q_6)\rho_6(t_6, a_6)da_6$$

Where $D_6$ is the mortality function and $n_6$ is the fertility function.

We assume following equations held

$$\left[\frac{dQ_6}{dt}\right]_6 = x_6\left[\frac{dQ_2}{dt}\right]_6 - x_2\left[\frac{dQ_6}{dt}\right]_6$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_6 - x_4\left[\frac{dQ_2}{dt}\right]_6$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_6 - x_6\left[\frac{dQ_3}{dt}\right]_6$$

Where $x_6, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_6$ and $g_6(P_6)Q_6$

Now when you add the term $(b_6Q_6)$ with $g_6(P_6)(Q_6)$ we are infect see is that $(b_6Q_6)$ is credited incrementally part of $g_6(P_6)(Q_6)$ or instant nearly, when subtracted from $\frac{dQ_6}{dt}$ , it means that $g_6(P_6)(Q_6)$ is slowly eaten by $\frac{dQ_6}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_6 = \mu_6(a_6) + (f_6)_6(P_6)_6 + (f_2)_6(Q)_6 \quad (2)$$

$$n_6 = \beta_6(a_6)h_6(Q_6), \quad (3)$$

Where $\mu_6$the age- dependent debit rate whit is $\mu_6(a_6)$ 0 $\forall a > 0$; $(f_6)_6$ is the density- dependent death rate with $f_6(0) = 0$, $(f_6^6)_6(P_6)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_6)_6(P_6) = (k_6)_6$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_6(0) = 0, (f_2^6)_6(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_6(Q_6) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 6, h^6(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_6 + (\rho_a)_6 = -\mu_6(a_6) + (f_6)_6P_6) + \big((f_2)_6(Q_6)\big)(\rho_6)_6 \quad (4)$$

$$\rho_6(t_6, 0) = h_6(Q_6)\int_0^\infty \beta_6(a_6)\rho_6(t, a_6)da, \quad (5)$$

$$\frac{dQ_6}{dt}$$

$$= -b_6Q_6 + g_6(P_6)Q_6 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_6$ and $Q_6$ of (4-6) satisfy

$$0 \ \rho_6(t, a_6) \ \rho_6^*(t, a_6),$$

$$0 \leq Q_6(t) \leq Q_6^*(t),$$

Where $\rho_6^*(t, a_6)$ and $Q_6^*(t)$, are solutions of

$$(\rho_t^*)_6 + (\rho_a^*)_6 = -(\mu_6(a_6) + (f_6)_6(P_6^*) + (f_2)_6 Q_6^*))\,\rho_6^*, \quad (7)$$

$$\rho_6^*(t, 0) = \int_0^\infty \beta_6(a_6)\,\rho_6^*(t, a_6)\,da_6 \quad (8)$$

$$\frac{dQ_6^*}{dt} = -b_6\,Q_6^* + g_6(P_6^*)Q_6^*, \quad (9)$$

With $P_6^*(t) = \int_0^\infty \rho_6^*(t, a_6)da_6$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_6^*(t)}{dt} = P_6^*(t)\big(r_6 - (f_6)(P_6^*) - (f_2)_6(Q^*)\big), \quad (10)$$

$$\frac{dQ_6^*(t)}{dt} = Q_6^*(t)(-b_6 + g_6(P_6^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_6(a)e^{-\int_0^a \mu_6(\tau)dt}\,e^{-ra}\,da_6 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_6^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(t)dt}\, da_6 > 1, \qquad (H1)$$

Which guarantees $r_6 > 0$.

*Remark.* If $\beta_6(a) = [\beta_0]_6 a_6 e^{-\alpha_6 a_6}, \mu_6(a_6) \equiv \mu_6$ then the solution of (12) is

$$r_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_6]_6(P_6) = 0, [f_2](Q_6) = \lambda_6 Q_6\ h_6(Q_6) = 1, \qquad and\ \ g_6(P_6) = c_6 P_6$$

$$[\rho_6]_6 + [\rho_a]_6 = -(\mu_6 + \lambda_6 Q_6)\rho_6$$

$$\rho_6(t,0) = \int_0^\infty [\beta_0]_6 a_6 e^{-\alpha_6 a_6} \rho_6(t, a_6)\, da_6$$

$$\frac{dQ_6}{dt} = -b_6 Q_6 + c_6 P_6 Q_6$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_6}{dt} = \omega_6 P_6 - \lambda_6 P_6 Q_6$$

$$\frac{dQ_6}{dt} = b_6 Q_6 - c_6 P_6 Q_6$$

Where $\omega_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$

## 49. THE EQUILIBRIUM $(\rho_0)_6(a_6), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_6(a_6), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_6 + (\rho_a)_6 = -\big(\mu_6(a_6) + (f_6)_6(P_6)\big)\rho,$$

$$\rho_6(t,0) = \int_0^\infty \beta_6(a_6)\rho_6(t,a_6)da_6 \qquad (13)$$

The solution of (63) has the form of

$$\rho_6(t,a_6) = Ce_6^{r(t-a_6)} - \int_0^a (\tau)d\tau - \int_0^6 (f_6)_6\big(P_6(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_6(t) = C_6 e^{r_6 t - \int_0^t (f_6)_6(P_6(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_6^{-6}(r)$ if $k_6 > r$. This is equivalent to

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{(k_6)_6 a_6} da_6 < 1.$$

In fact, if we define

$$R_6(s_6) = \int_0^\infty \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau)d\tau} e^{-s_6 a_6}\, da_6$$

$R(s)$ is a decreasing function. $R_6(r_6) = 1$ And $R_6(k_6)_6 < 1$ imply $r_6 < (k_6)_6$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_6(a_6), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{k_6 a_6} da_6$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_6 \equiv \int_0^\infty (\rho_0)_{66}(a_6)da_6$ *satisfies*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{-(f_6)_6(P_0)a_6} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_6(a_6), 0)$

Let $u_6(t, a_6) = \rho(t, a_6) - \rho_0(a_6)$ Then

$$N_6(t) \equiv \int_0^\infty u_6(t, a_6)da_6 = \int_0^\infty \rho_6(t, a_6)da_6 - \int_0^\infty (\rho_0)_6(a_6)da = P_6(t) - P_0)_6$$

And

$$(u_6)_6 + (u_a)_6 = (\rho_t)_6 + (\rho_a)_6 - (\rho_0^6)_6(a_6)$$

$$= -\mu_6(a_6) + (f_6)_6(P_6) + (f_2)_6(Q_6))\,(u_6 + (\rho_0)_6) + \mu_6(a_6)(\rho_0)_6 + (f_6)_6(P_0)_6(\rho_0)_6$$

$$= -\mu_6(a_6)u_6 - (f_6)_6(P_6) + (f_2)_6(Q_6))\,u_6 - (\rho_0)_6\,(f_6)_6(P_6) - (f_6)_6(P_0)_6$$

$$+ (f_2)_6(Q_6)), \qquad (15)$$

Where $(\rho_0^6)_6(a_6)$ denotes $d(\rho_0)_6(a_6)/da_6$

Expansion $(f_6)_6(P_6)$ and $(f_2)_6(Q_6)$ at $(P_0)_6$ and 0 yield respectively

$$(f_6)_6(P_6) = (f_6)_6((P_0)_6 + N_6)$$

$$= (f_6)_6(P_0)_6 + (f_6^6)_6(P_0)_6 N_6 + o_6(N_6), \qquad (16)$$

$$(f_2)_6(Q_6)$$

$$= (f_2^6)_6(0)Q + o_6(Q_6) \quad (17)$$

Write $f_6(P_0) = f_{60}, f_6'(P_0) = f_{60}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_6)_6 + (u_a)_6 \approx -(\mu_6(a_6) + (f_{60})_6)u - (f_{60}^6)_6(\rho_0)_6(N_6) - (f_{20}^6)_6(\rho_0)_6(Q_6)$$

$$u_6(t,0) = \rho_6(t,0)(\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a) + (\rho_0)_6(a_6)da_6 - (\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a_6)da_6 + (h_6(Q_6) - 1)(\rho_0)_6(0) \quad (18)$$

Expanding

$$h_6(Q_6) = 1 + (h^6)_6(0)Q_6 + 0(Q_6) \equiv 1 + h_0'Q_6 + 0Q_6$$

And substituting it into (18), we arrive at

$$u_6(t,0) \approx \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + (h_{a_6}')(\rho_0)_6Q_6$$

The linearization is the following :

$$[u_6] + [u_6] = -(\mu_6(a_6) + [t_{60}]_6u_6 - [\alpha_6]_6(a_6)N_6 - [\alpha_2]_6(a_6)Q_6$$

$$u_6(t,0) = \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + [\xi_6]Q_6 \quad (19)$$

$$\frac{dQ_6}{dt} = -b_6Q_6$$

where

$$[\alpha_6]_6(a_6) = [t_6^6 a]_6[p, o]_6(a_6)d2(a_6) = [f_{20}^6]_6(\rho\, o\,)_6(a_6)\xi = [\mathrm{h\ o}^6]_6[\rho \mathrm{o}]_6(\mathrm{o}) \text{ and } \mathrm{b}_6 = \mathrm{b}_6 - g_6(P_6)$$

If $g_6(P_0) > b$, clearly, $(\rho_0)_6(a_6), 0)$ is unstable and the predator population is persistent. Assume $g_6(P_0)_6 < b$.

Then, $Q_6(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_6 + u_a = -(\mu_6(a_6) + (f_{60})_6)u_6 - (\alpha_6)_6(a_6)N_6$$

$$= -(\mu_6(a_6) + f_{60})u_6 - \alpha_6(a_6)\int_0^\infty u_6(t, a_6)da_6$$

$$u_6(t, 0) = \int_0^\infty \beta_6\,(a_6)u_6(t, a_6)da_6$$

Using the method developed by Cushing [2, 3], we suppose

$$u_6(t, a_6) = \tilde{\mathrm{u}}_6(a_6)^{e^{\gamma(t-a_6)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_6}\, da_6$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_6)$ gives,

$$\tilde{\mathrm{u}}_6(a_6) = e^{-\int_0^a \mu_6(\tau_6)d\tau - (f_{60})_6 a}\left(\tilde{\mathrm{u}}_6(0) - w\int_0^a e^{\int_0^\tau \mu_6(s_6)ds + (f_{60})}\tau_6\,(\alpha_6)_6(\tau_6)e^{\gamma_6\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_6(a_6)e^{-\gamma_6 a_6 - (f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6)d\tau}\right)\tilde{\mathrm{u}}_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a(\alpha_6)_6(\tau_6)e^{-\gamma_6\tau_6}e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds}d\tau_6\,da_6$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}da\right)(\tilde{u}_6(0))_6$$

$$+\left(-6-\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds+y\tau}\alpha_6(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(t_6)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_6\,(a_6)Y_6(\gamma_6,a_6)da_6\right)u_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6=0,$$

(24)

$$\left(\int_0^{\infty}Y_6(\gamma_6,a_6)da_6\right)\tilde{u}_6(0)$$

$$+\left(-1-\int_0^{\infty}Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6\right)w=0.$$

Equqtion (20) has a solution of the form of $u_6(t,a_6)=u_6(a_6)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_6(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) Y^{-6}(\gamma_6, a_6) d\tau_6$. We obtain the characteristic equation

$$f_{60}^{l}{}_6 \left( \int_0^\infty Y_6(\gamma_6, a_6) da_6 \right) \left( \int_0^\infty \beta_6(a_6) Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$+ \left( 1 - \int_0^\infty \beta_6(a_6) Y_6(\gamma_6, a_6) da \right) \left( 1 + (f_{60}^{l})_6 \int_0^\infty Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$= 0. \qquad (25)$$

Since

$$(\rho_0)_6(a_6) = (\rho_0)_6(0) e^{-\int_0^a \mu_6(\tau_6) d\tau_6 - (f_{60})_6 a_6}$$

$$I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) e^{\gamma_6 \tau_6 + (f_{60})_6 \tau + \int_0^\tau \mu_6(s_6) ds_6} d\tau$$

$$= \frac{(\rho_0)_6(0)}{\gamma_6} (e^{\gamma_6 a_6} - 1)$$

Let

$$E_6(a_6) = e^{-(f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6) d\tau_6}$$

This relationship is cardinal one. It is useful in quantum computation and other black hole dynamics problems

**Quantum space and quantum time**

(2716) **Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.**

(2717) **In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.**

(2718) **They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.**

(2719) **When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.**

(2720) **Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.**

## The Planck scale

(2721) **Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".**

(2722) **Here ħ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.**

(2723) **The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.**

(2724) **Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)**

(2725) **Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.**

## Physical applications

**Some phenomena where quantum gravity plays a major role are the following.**

### The microscopic structure of spacetime

(2726) **If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.**

(2727) **Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length**

(2728) **Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.**

(2729) **A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.**

**Early cosmology**

(2730) **According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.**

(2731) **However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.**

(2732) **A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.**

(2733) **Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".**

(2734) **One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.**

**Black holes**

(2735) **Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.**

(2736) **First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.**

(2737) **Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.**

(2738) **What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.**

(2739) **Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.**

**Astrophysical effects**

(2740) **Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.**

(2741) **An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.**

(2742) **The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.**

## Main approaches

(2743) **The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.**

(2744) **Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.**

(2745) **The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.**

## String theory

(2746) **String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.**

(2747) **Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.**

(2748) **String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.**

(2749) **Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.**

(2750) **However, when summed the entire series appears to be (=) divergent.**

(2751) **In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.**

(2752) **Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.**

(2753) **Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.**

(2754) **In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.**

## Loop quantum gravity and spin foams

(2755) **Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.**

(2756) **The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity.**

(2757) **Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime.**

(2758) **The granular properties of space can be (=) explicitly computed.**

(2759) **In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory.**

(2760) **The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks.**

(2761) **Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure.**

(2762) **The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit.**

(2763) **Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism.**

(2764) **This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism.**

## Noncommutative geometry

(2765) **Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime.**

(2766) **Quantum mechanics teaches us that dynamical quantities are (=) noncommutative -- meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed.**

(2767) **It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry.**

(2768) **A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces.**

(2769) **Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces.**

(2770) **This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity.**

**Others**

(2771) **Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays**

(2772) **. Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection.**

(2773) **Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible.**

(2774) **The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.**

**Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016**

(2775) **These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field.**

(2776) **The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets.**

(2777) **Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state.**

(2778) **Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands.**

(2779) **The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI**

(2780) **The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9].**

(2781) **This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive.**

(2782) **More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.**

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

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| (2783) | **Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.** |
| (2784) | **Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.** |
| (2785) | **The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.** |
| (2786) | **This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.** |
| (2787) | **The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.** |
| (2788) | **This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.** |
| (2789) | **For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.** |
| (2790) | **The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society** |

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

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| (2791) | **Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.** |
| (2792) | **We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.** |
| (2793) | **When this is true—as is often the case—the system serves as (=) a "bath" for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called thermalization.** |
| (2794) | **However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.** |
| (2795) | **Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1).** |
| (2796) | **Bringing them to thermal equilibrium requires (e) an external bath.** |

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| (2797) **But so far, there has been no rigorous proof of (e) the existence of such many-body localization. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption.** |
| (2798) **The concept of localization was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations).** |
| (2799) **The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century.** |
| (2800) **But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information.** |
| (2801) **Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment.** |
| (2802) **For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood.** |
| (2803) **It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7].** |
| (2804) **Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature.** |
| (2805) **One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.** |
| (2806) **However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].** |
| (2807) **Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.** |
| (2808) **Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.** |
| (2809) **As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked.** |
| (2810) **His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice.** |
| (2811) **An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site.** |
| (2812) **This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself).** |

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| (2813) **This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition.** |
| (2814) **However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime:** |
| (2815) **These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects.** |
| (2816) **It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies).** |
| (2817) **This is a realistic scenario for (e) all known systems.** |
| (2818) **On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects.** |
| (2819) **On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other.** |
| (2820) **No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true.** |
| (2821) **His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found.** |
| (2822) **Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization?** |
| (2823) **Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects?** |
| (2824) **Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions.** |
| (2825) **So in those systems, the quantum phase transition does indeed exist.** |
| (2826) **But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension.** |
| (2827) **Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters.** |

**Focus: Biological Cells Form Electric Circuits July 1, 2016• Physics 9, 75 Cells that are electrically active and that also produce light for easy voltage monitoring could lead to new studies of heart arrhythmias and possibly bio-computing. Figure caption expand figure H. M. McNamara et al., Phys. Rev. X (2016) Cell circuit. Waves of electrical activity propagate in a loop within a collection of cells engineered with electrical and optical properties (see video below).Some sentences, pictures and videos are deleted due to constraints on space.**

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| (2828) **The human heartbeat is produced by (e) electrical pulses that propagate through (e&eb) cardiac tissue, causing (eb) rhythmic muscle contraction.** |
| (2829) **Researchers have previously engineered cells to create (eb) an artificial tissue capable of producing (eb) coordinated electrical activity, and now a team has added the ability to monitor (e&eb) their electrical state by (e) detecting fluorescent emission.** |
| (2830) **They have also fashioned the cells into (e&eb) "living circuits" that might act as (=) model systems for studying heart behavior.** |
| (2831) **Constructing electrically active tissues from their basic building blocks by (e) genetically engineering cells allows (eb) researchers to study a "clean" system similar to (e) the more complicated cardiac tissue.** |
| (2832) **These systems let them investigate, for example, the electrical signals that lead to (eb) the potentially fatal irregular heartbeat called arrhythmia.** |
| (2833) **Previously researchers have introduced (eb) ion channels—proteins that conduct ions through (e&eb) cell membranes—into human kidney cells to make (eb) electrically "excitable" tissues [1].** |
| (2834) **Adam E. Cohen of Harvard University and his co-workers decided to add more ingredients to such cells, so that the electrical activity could be (=) light-activated and easily visualized.** |
| (2835) **They incorporated four additional gene products into (e&eb) human embryonic kidney (HEK) cells.** |
| (2836) **Two were ordinary ion channels; the other two were (=) a blue-light-activated ion channel (a light-sensitive switch) and (e&eb) a protein that fluoresces with (e&eb) an intensity that depends on (=) the voltage across the membrane in which it is embedded.** |
| (2837) **Light activation of electrically excitable cells has been shown before [2]. But Cohen and colleagues have now included a convenient "readout" mechanism (the fluorescent protein) and have used patterning methods to shape (e&eb) the cell cultures into "tissue circuits."** |
| (2838) **They grew rings of cells a few millimeters in diameter that could act as (e&eb) oscillators that support (eb) circulating waves of electrical excitation.** |
| (2839) **These waves arise from (e) the collective behavior of the cells: a pulse produced by (e) one cell can trigger another from (e) its neighbors, eventually producing (eb) a coherent wave of excitation traveling through (e&eb) the tissue.** |
| (2840) **It's this coordination, in a repeating cycle of excitation and quiescence, which gives rise to (eb) rhythmic activity in cardiac tissue.** |

**H. M. McNamara et al., Phys. Rev. X (2016)**

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| (2841) **A pattern of blue light (not visible) illuminates these cells, defining (eb) a rectangular loop around which electrical waves can propagate.** |
| (2842) **At about 0:07 the rectangle begins to shrink in (eb) the vertical direction, causing t(eb) he cycle time to gradually decrease. (In other experiments, the cells were grown in (eb) fixed, ring-shaped patterns.)** |
| (2843) **The ring of engineered HEK cells acted like an electronic "ring oscillator," a device with a characteristic frequency to (e&eb) its circulating electrical signal.** |
| (2844) **The electrical waves in a ring oscillator can propagate (e&eb) in either direction around the ring, enabling (eb) it to serve as a binary memory unit.** |

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| (2845) **Cohen and colleagues found that in their HEK tissues, these circulating waves remained (eb) stable for more than two hours, with (e&eb) the pulses completing a few circuits every second.** |
| (2846) **The electrical signals were easily visualized as (=) waves of fluorescent emission.** |
| (2847) **The researchers say that these tissues could be used to (e) understand and predict the behavior of real cardiac cells, especially arrhythmias.** |
| (2848) **"In real heart, if you pace the tissue fast enough, *(eb) it starts skipping beats and eventually becomes (=) arrhythmic," says Cohen.** |
| (2849) **This state is a well-known precursor to (e) sudden cardiac death and is associated with (e&eb) the breakup of the usual regular waves of electrical excitation.** |
| (2850) **Some drugs also have the dangerous side effect of inducing (eb) arrhythmia, and Cohen says that the new system might offer a quick and convenient way of screening for such hazards. The work "is really very elegant and deserves attention," says physiologist Gil Bub of Oxford University in the UK. "It opens new research questions for people like me who are interested in wave dynamics in excitable systems."** |
| (2851) **There might also be practical uses for information processing in living tissues. "By combining patterns of cell growth and (e&eb) patterns of ion channel expression, we think we can build diodes, amplifiers, and other simple logic elements," says Cohen.** |
| (2852) **Andrew Adamatzky of the University of the West of England in Bristol, a specialist in "unconventional computing," foresees circuits that perform computations through (e&eb) the interactions of electrical waves. "This research might lead to fabrication of implantable bio-computers," he says. This research is published in Physical Review X. –Philip Ball Philip Ball is a freelance science writer in London and author of Invisible: The Dangerous Allure of the Unseen (University of Chicago Press, 2014).** |

**Optically Controlled Oscillators in an Engineered Bioelectric Tissue Harold M. McNamara, Hongkang Zhang, Christopher A. Werley, and Adam E. Cohen Phys. Rev. X 6, 031001 – Published 1 July 2016**

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| (2853) **Complex electrical dynamics in excitable tissues occur throughout (e&eb) biology, but the roles of individual ion channels can be difficult to determine due to (e) the complex nonlinear interactions in (eb) native tissue.** |
| (2854) **Here, authors ask whether they can engineer a tissue capable of (e) basic information storage and processing, where (e) all functional components are known and well understood.** |
| (2855) **Authors develop a cell line with (e&eb) four transgenic components: two to enable (eb) collective propagation of electrical waves and (e&eb) two to enable optical perturbation and (e&eb) optical readout of membrane potential.** |
| (2856) **They pattern the cell growth to define (eb) simple cellular ring oscillators that run stably for >2 h (∼104 cycles) and that can store data encoded in (eb) the direction of electrical circulation.** |
| (2857) **Using (e) patterned optogenetic stimulation authors probe the biophysical attributes of (e) this synthetic excitable tissue in detail, including (e) dispersion relations, (e&eb) curvature-dependent wave front propagation, (e&eb) electrotonic coupling, and (e&eb) boundary effects.** |
| (2858) **Authors then apply the biophysical characterization to (e) develop an optically reconfigurable bioelectric oscillator.** |
| (2859) **These results demonstrate the feasibility of engineering bioelectric tissues capable of (e&eb) complex information processing with optical input and output. Received 12 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevX.6.031001 This article is available under the terms of the** |

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**Scientists detect radio emission from a nearby brown dwarf July 6, 2016 by Tomasz Nowakowski report**

(2860) **K2 and Spitzer images of W0607+24: The position of W0607+24 is marked by (e) a red circle. SDSS J060738.06+242939.5, marked by (e) an orange circle, are the brightest star in the K2 image.**

(2861) **The three faint (r = 18.3 − 19.6) stars marked by cyan circles, contribute significant flux to (e) the K2 measurement but are (-) negligible at Spitzer's longer wavelength and (e &êb) higher resolution. Other stars with r < 20 in SDSS are marked with blue circles. Credit: Gizis et al., 2016.**

(2862) **(Phys.org)—Located some 23 light years away, the brown dwarf designated WISEP J060738.65+242953.4 or W0607+24 for short, turns out to be (=) a source of radio emission.**

(2863) **According to a research paper published July 4, this substellar object showcases (eb) quiescent radio emission, making it one of (e) the most radio-faint, ultra-cool dwarfs yet detected.**

(2864) **Brown dwarfs like W0607+24 are objects that are too large to be called planets and too small to be stars. With a mass below that necessary to maintain hydrogen-burning nuclear fusion reactions, they are much cooler and dimmer than (eb) main sequence stars.**

(2865) **These objects could be also sources of (e) radio emission, but the cause of this process is still not completely understood.**

(2866) **Due to its proximity, W0607+24 is (=) an excellent target for studying the nature of (e) brown dwarfs.**

(2867) **It is the nearest known late-L dwarf in the northern hemisphere, and the third-nearest in (=) the whole sky, (eb) allowing researchers to study phenomena occurring in brown dwarfs.**

(2868) **That's why a team of astronomers, led by John Gizis of the University of (e)Delaware, Newark, carried out a series of mid-infrared, radio, and optical observations of W0607+24 to study any activity in this curious substellar object.**

(2869) **Motivated by the results from NASA's prolonged Kepler mission, named K2, which obtained (eb) long time-series photometry of W0607+24; they observed this object simultaneously with the Karl G. Jansky Very Large Array (VLA) in New Mexico and NASA's Spitzer Space Telescope.**

(2870) **The most important results were provided by VLA, as the team found that W0607+24 is (=) a source of quiescent, non-bursting, radio emission.**

(2871) **The scientists noted that the discovery of this process in this nearby object could be (=) crucial for our understanding of this process in (eb) ultracool dwarfs in general.**

(2872) **"W0607+24 is a valuable test case for (e) understanding the mechanisms of radio emission in ultracool dwarfs.**

(2873) **Some of the radio-detected objects emit (eb) rapid, intense bursts with high circular polarization, often found to occur periodically and attributed to coherent emission due to (e) the electron cyclotron maser instability," the paper reads.**

(2874) **The astronomers also suggested that the relative radio faintness of W0607+24 may be due to (e) its orientation if (e) the emission is concentrated at the poles.**

(2875) **However, follow-up observations are needed to better characterize (eb) the detected radio emission and to eliminate (e) all the uncertainties.**

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| (2876) **The team also hopes for more significant findings when (e) studying this nearby substellar object in the future.** |
| (2877) **"As the nearest known northern hemisphere late-L dwarf, this offers many possibilities for future studies. For example, it could serve as (=) an useful standard for (e) line-broadening and (e&eb) polarization studies.** |
| (2878) **If there is a planetary system aligned with (e&eb) (the primary's rotation axis, then we would expect very little radial velocity signal but a strong astrometric signal," the researchers concluded.** |
| (2879) **Moreover, the team plans to monitor (e&eb) additional late-M and L dwarfs in other K2 campaigns, expecting even more promising results.** |
| (2880) **They are more distant than W0607+24, but because they are warmer, (eb) some are brighter and higher signal-to-noise than that described in this paper.** |
| (2881) **In addition to detecting the radio emission, (eb)the scientists also estimated the age, mass and radius of W0607+24** |
| (2882) **. According to the study, the object is less than two billion years old, has (e) a maximum mass of about 0.055 solar masses and a radius approximately of 0.1 solar radii.** |
| (2883) **More information: WISEP J060738.65+242953.4: A Nearby. Pole-On L8 Brown Dwarf with Radio Emission, arXiv: 1607.00943 [astro-ph.SR] arxiv.org/abs/1607.00943** |
| (2884) **Authors present a simultaneous, multi-wavelength campaign targeting the nearby (7.2 pc) L8/L9 (optical/near-infrared) dwarf WISEP J060738.65+242953.4 in (eb) the mid-infrared, radio, and optical.** |
| (2885) **Spitzer Space Telescope observations show no (e) variability at the 0.2% level over 10 hours each in the 3.6 and 4.5 micron bands.** |
| (2886) **Kepler K2 monitoring over 36 days in Campaign 0 rules out (e) stable periodic signals in (eb) the optical with amplitudes great than 1.5% and periods between 1.5 hours and 2 days.** |
| (2887) **Non-simultaneous Gemini optical spectroscopy detects lithium, constraining (e) this L dwarf to be less than ~2 Gyr old, but no Balmer emission is observed.** |
| (2888) **The low measured projected rotation velocity (v sin i < 6 km/s) and lack of variability are (=) very unusual compared to other brown dwarfs, and authors argue that this substellar object is (=) likely viewed pole-on.** |
| (2889) **Authors detect quiescent (non-bursting) radio emission with (e&eb) the VLA.** |
| (2890) **Amongst radio detected L and T dwarfs, it has (e) the lowest observed L_nu and the lowest v sin i.** |
| (2891) **They discuss the implications of pole-on detection for (e) various proposed radio emission scenarios. © 2016 Phys.org** |

**A new look at the galaxy-shaping power of black holes July 6, 2016**

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| (2892) **It wasn't until the advent of X-ray astronomy that the full picture of (e) the galaxy-shaping power of black holes began to emerge (eb) with the ability to see plasma.** |
| (2893) **Data from a now-defunct X-ray satellite is providing (eb) new insights into the complex tug-of-war between galaxies, the hot plasma that surrounds them, and the giant black holes that lurk in their centres.** |

(2894) **Launched from Japan on February 17, 2016, the Japanese space agency (JAXA) Hitomi X-ray Observatory functioned for just over a month before contact was lost and (e&eb) the craft disintegrated.**

(2895) **But the data obtained during those few weeks was enough to paint (eb) a startling new picture of the dynamic forces at work within galaxies.**

(2896) **New research, published in the journal Nature today, reveals data that shows just how important the giant black holes in galactic centres are to the evolution (eb) of the galaxies as a whole.**

(2897) **"We think that supermassive black holes act(e&eb) like thermostats," said Brian McNamara, University Research Chair in Astrophysics at the University of Waterloo. "They regulate (e&eb) the growth of galaxies."**

**Champagne bubbles of plasma**

(2898) **During its brief life, the Hitomi satellite collected X-ray data from (e) the core of the Perseus cluster, an enormous gravitationally-bound grouping of (e) hundreds of galaxies.**

(2899) **Located some 240 million light years from earth, the Perseus cluster is one of (e) the largest known structures in the universe.**

(2900) **The cluster includes not only the ordinary matter that makes up the galaxies, but an (e) "atmosphere" of hot plasma with a temperature of tens of millions of degrees, as well as a halo of (e) invisible dark matter.**

(2901) **Earlier studies, going back to the 1960s, have shown that each of the galaxies in the cluster - and indeed most galaxies - likely contains (e) a supermassive black hole in its centre, an object 100 million to more than ten billion times as massive as our sun**

**New look at the galaxy-shaping power of black holes**

(2902) **Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving.**

(2903) **"These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science.**

(2904) **"Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy."**

(2905) **This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return.**

(2906) **The remaining matter gets absorbed into the black hole, adding to (e) its mass.**

(2907) **The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne.**

(2908) **The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution.**

(2909) **Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The**

**answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form.**

(2910) **"Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara.**

(2911) **"It creates (eb) these bubbles, and the bubbles keep the plasma hot.**

(2912) **That's what prevents galaxies from (e) becoming even bigger than they are now."**

(2913) **Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.**

(2914) **In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.**

(2915) **Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.**

(2916) **"It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.**

(2917) **"It's like nature's thermostat, which keeps these galaxies from (e) growing.**

(2918) **If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."**

**A new look at the galaxy-shaping power of black holes**

(2919) **An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the**

(2920) **McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.**

(2921) **Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.**

**Doomed satellite missions**

(2922) **Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.**

(2923) **At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.**

(2924) **Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a**

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| **leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.** |
| (2925) **The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.** |
| (2926) **These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.** |
| (2927) **"But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people** |
| (2928) **. It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."** |
| (2929) **The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.** |
| (2930) **"This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website** |

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

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| (2931) **Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.** |
| (2932) **They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.** |
| (2933) **However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.** |
| (2934) **Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.** |
| (2935) **X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well.** |
| (2936) **The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma.** |
| (2937) **These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6.** |
| (2938) **Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus.** |

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| (2939) | **A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate.** |
| (2940) | **Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure.** |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (2941) | **This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy.** |
| (2942) | **Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole.** |
| (2943) | **But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy.** |
| (2944) | **A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole.** |
| (2945) | **It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth.** |
| (2946) | **"There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA).** |
| (2947) | **This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies.** |
| (2948) | **However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated.** |
| (2949) | **In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1.** |
| (2950) | **We know dark matter exists only from (e) its gravitational effects.** |
| (2951) | **It holds together galaxies and (e&eb) galaxy clusters.** |
| (2952) | **Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years.** |
| (2953) | **To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies.** |
| (2954) | **They used star motions as tracers to (e) weigh the galaxies' central black holes.** |
| (2955) | **X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto.** |
| (2956) | **They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone.** |
| (2957) | **This connection is likely to be related to (e&eb) how elliptical galaxies grow.** |

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| (2958) **An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together.** |
| (2959) **Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole.** |
| (2960) **"In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp** |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (2961) **The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the …** |
| (2962) **Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness?** |
| (2963) **Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness.** |
| (2964) **"Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain."** |
| (2965) **Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness.** |
| (2966) **"With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication."** |
| (2967) **Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain** |

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| **changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories.** |
| (2968) **The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen.** |
| (2969) **In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer.** |
| (2970) **Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants.** |
| (2971) **These were treated as (=) "aware" and "unaware" trials, respectively.** |
| (2972) **Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it.** |
| (2973) **The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another.** |
| (2974) **Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion.** |
| (2975) **No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness.** |
| (2976) **"We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said.** |
| (2977) **"Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness."** |
| (2978) **The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner.** |
| (2979) **Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain.** |
| (2980) **"We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said.** |
| (2981) **"This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients** |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

(2982) **Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter.**

(2983) **As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined.**

(2984) **Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions.**

(2985) **"Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.**

(2986) **The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.**

(2987) **However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.**

(2988) **The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.**

(2989) **Noda and his colleagues instead studied the change in emissions over (e&eb) time.**

(2990) **By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.**

(2991) **Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).**

(2992) **The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.**

(2993) **Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.**

(2994) **"The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp**

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(2995) **The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking "palm tree" organization.**

(2996) **The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society**

(2997) **(Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible?**

(2998) **With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem.**

(2999) **This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other.**

(3000) **According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible.**

(3001) **This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges.**

(3002) **In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century.**

(3003) **Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching.**

(3004) **In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom.**

(3005) **Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth.**

**Spheres on trees**

(3006) **Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration.**

(3007) **They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement.**

(3008) **Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum.**

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| (3009) **By visualizing the problem in this way, the researchers noticed that these particular graphs don't (e) have lots of branches extending from lots of (e) other branches.** |
| (3010) **Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure.** |
| (3011) **The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network.** |
| (3012) **As a result, it provides (eb) an important clue for finding the global minimum.** |
| (3013) **It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum.** |
| (3014) **This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large.** |

**Thompson's problem 2**

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| (3015) **By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society** |
| (3016) **The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima.** |
| (3017) **The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node.** |

**Implications for Thomson's and Smale's problems**

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| (3018) **Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before.** |
| (3019) **"Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around."** |
| (3020) **As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem.** |
| (3021) **"I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view."** |
| (3022) **In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges.** |
| (3023) **From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear.** |

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| (3024) | **"While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK.** |
| (3025) | **"For larger numbers, there will likely be (=) multiple funnels.** |
| (3026) | **Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum."** |
| (3027) | **Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.** |
| (3028) | **"The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.** |
| (3029) | **"Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org** |

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**Introduction:**

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price, University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss.**

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| **(481)** | **Black Holes are regions of (e) space in which gravitational fields are so strong that no particle or signal can escape (eb) the pull of gravity.** |
| **(482)** | **The boundary of this no-escape region is (=) called the event horizon, since distant observers outside the black hole cannot (e) see (cannot get light from) events inside.** |
| **(483)** | **Although the fundamental possibility of such an object exists within (eb) Newton's classical theory of gravitation, Einstein's theory of gravity makes (eb)black holes inevitable under some circumstances.** |
| **(484)** | **Prior to the early 1960s, black holes seemed to be (=) only an interesting theoretical concept with no (e) astrophysical plausibility, but with (e&eb) the discovery of quasars in 1963 it became clear that very exotic astrophysical objects could exist.** |
| **(485)** | **Nowadays it is taken for granted that black holes do exist in (eb) at least two different forms.** |
| **(486)** | **Stellar mass black holes are (=) the endpoint of the death of some stars, and supermassive black holes are the result of (e) coalescences in the centers of most galaxies, including our own.** |
| **(487)** | **No signal can propagate from (e) inside a black hole, but the gravitational influence of a black hole is (=) always present.** |

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| **(488)** | **This influence does not (e) propagate out of the hole; it is permanently present outside, and depends only on (=) the total amount of mass, angular momentum, and electric charge that have gone into forming (eb) the hole.** |
| **(489)** | **Black holes can be detected through the influence of this (e&eb) strong gravity on the surroundings just outside the hole.** |
| **(490)** | **In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them.** |
| **(491)** | **Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors.** |

**Classical vs. relativistic black holes**

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| **(492)** | **Something like a black hole exists within (eb) Newton's classical theory of gravity.** |
| **(493)** | **In that theory, an energy argument tells us that there is (=) an escape velocity $v_{esc}=\sqrt{2GM/R}$ from (e) the surface of any spherical object of (e) mass M and radius R** |
| **(494)** | **If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity.** |
| **(495)** | **Thus the condition for such an "unseeable" object is $R<2GM/c^2$** |
| **(496)** | **In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape.** |
| **(497)** | **This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity.** |
| **(498)** | **To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it.** |
| **(499)** | **The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c .** |
| **(500)** | **This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed.** |
| **(501)** | **An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity.** |
| **(502)** | **Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable.** |

**Stationary black holes**

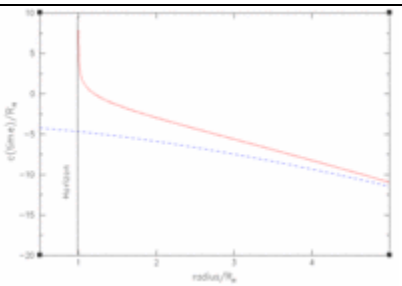
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| **(503)** | **In the Newtonian theory, gravity is described by (e) the potential $\Phi$.** |
| **(504)** | **Inside a spherical object the form of (e) $\Phi(r)$ depends on (=) the interior structure, but in the vacuum outside (e) matter the potential $-GM/r$ depends only on (=) the interior mass.** |
| **(505)** | **Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object.** |

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| **(506) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward.** |
| **(507) This horizon radius RH=2GM/c2 is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory.** |
| **(508) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution.** |
| **(509) Radius RHhere actually means that the area of the event horizon is (=) 4πR2H.** |
| **(510) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution.** |
| **(511) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity.** |
| **(512) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own.** |
| **(513) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole.** |
| **(514) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density.** |
| **(515) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed.** |
| **(516) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole.** |
| **(517) In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter.** |
| **(518) Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum.** |
| **(519) In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior.** |
| **(520) The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object.** |
| **(521) Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum.** |
| **(522) These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.]** |
| **(523) All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time.** |
| **(524) But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?"** |
| **(525) The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point.** |

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| **(526) One of these coordinates is (=) called the "coordinate time."** |
| **(527) Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate.** |



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| **(528) Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole** |
| **(529) At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole.** |
| **(530) Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations.** |
| **(531) For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree** |
| **(532) An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events.** |
| **(533) The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving.** |
| **(534) Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle.** |
| **(535) The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return.** |
| **(536) But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon.** |
| **(537) Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future.** |
| **(538) The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?"** |

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| **Black hole parameters** |

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| **(539) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin).** |
| **(540) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun.** |
| **(541) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits.** |
| **(542) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form.** |
| **(543) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds.** |
| **(544) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2.** |
| **(545) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed.** |
| **(546) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer.** |
| **(547) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2)** |
| **(548) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials.** |
| **(549) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable** |
| **(550) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite.** |
| **(551) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole.** |
| **(552) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars.** |
| **(553) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity.** |
| **(554) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity.** |
| **(555) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite.** |
| **(556) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity.** |
| **(557) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum.** |

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| **(558) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship.** |
| **(559) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= (4πG2/c4)[(M+M2−a2−−−−−−−√)2+a2].** |
| **(560) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored.** |

**Dynamical black holes**

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| **(561) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape.** |
| **(562) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, a horizon can be born.** |
| **(563) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid.** |
| **(564) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward.** |
| **(565) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations.** |
| **(566) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes.** |
| **(567) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3.** |
| **(568) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/** |
| **(569) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers.** |
| **(570) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission.** |
| **(571) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.** |
| **(572) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.** |
| **(573) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.** |
| **(574) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.** |

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| **(575) The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature 6×10−8 (Mo/M) K.** |
| **(576) Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.** |

Then $[Y_6]_6[\gamma_6]_6, a_6) = E_6(a_6)e^{-\gamma_6 a_6}$. by a simple calculation,

$$\int_0^\infty \beta_6(a_6)Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma}\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}\ da\right);$$

$$\int_0^\infty Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma_6}\left(\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}\ da_6\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_6(0)}{\gamma_6}[f_{60}^6]\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}da_6\ (\int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6)$$

$$+\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right).$$

$$\left(1 + \frac{(\rho_0)_6(0)}{\gamma_6}(f_{60}^{'})_6\left(\int_0^\infty E_6(a_6)da_6 - \int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}da_6\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right)(1 + \frac{(f_{60}^{'})_6}{\gamma_6}(\rho_0)_6(0)\int_0^\infty E_6(a_6)da_6) = 0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1 - \int_0^\infty \beta_6(a_6) Y_6(\gamma_6 a_6) da_6 = 1 - \int_0^\infty \beta_6(a_6) E_6(a_6) e^{-\gamma_6 a_6} da_6 \neq 0,$$

Since $\int_0^\infty \beta_6(a_6) Y_6(\gamma_6 a_6) \int_0^a \alpha_6(\tau_6) Y^{-6}(\gamma_6, \tau_6) d\tau_6 da_6 > 0$. Therefore, the characteristic solution is

$$\gamma_6 = -(f_{60}')_6 (\rho_0)_6 \int_0^\infty E_6(a_6) da_6.$$

Since $f_6'(P_0)_6 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_6), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE OF DEBIT TYPES AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_6), 0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_6 + (\mu_6(a_6) + [f_6]_6(P_6^*) + [f_2]_6(Q_6^*))P_6^* = 0, \quad (26)$$

$$p_6^*(0) = h_6(Q_6^*) \int_0^\infty \beta_6(a_6) p_6^*(a_J) p_6^*(a_6) da_6 \quad (27)$$

$g_6(P_6^*) - b_6 = 0,$ (28)

Where $P_6^*$=$\int_0^\infty p_6^*(a_6)\, da_6$

Equation (28) yields $P^* = g^{-6}(b)$,

From (26),

$p_6^*(a) = p_6^*(0)e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6}$

Substituting into (27), we need to solve

$$\left(1 - h_6(Q_6^*)\int_0^\infty \beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6} da_6\right) p_6^*(0) = 0$$

For $Q_6^*$

Fix $p_6^*(0)$ and $P_6^*$ Let

$$H_6(z_6) = 1 - h_6(z_6)\int_0^\infty \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6} da_{,6}$$

And assume that

$$[k_3]_6 \int_0^\infty \beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(g_6^{-6}(b_6))a_6 - [k_2]_6 a_6}\, da_6 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_6(z_6)] > 0.$$

Clearly,

$$\frac{dH_6}{dz_6} = h_6(z)([f_2^{'}]_6(z_6))\int_0^{\infty} a_6\beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - f_6(P_6^*)a_6 - f_2(z_6)a_6}\, da_6$$

$-h^{'}(z_6)\int_0^{\infty}\beta_6(a_6)\, e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6}\quad da_6$

By the assumptions on $\left[f_2^6\right]_6 and\ [h^6]^6 dH_6/dz > 0$. Hence, if

$$H_6(0) = 1 - \int_0^{\infty}\beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6}da_6 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_6(0) < 0$ follows if $g_6^{-}(b_6) < [f_6^{-6}]_6(r_6)$, where $r_6$ is the solution of (12).

Indeed, since

$$R_6(s_6) = \int_0^{\infty}\beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 e^{-y_6^a}}da_6$$

Is a decreasing function of $s_6$ when $P_6^* = [g^{-6}]_6(b_6) < [f_6^{-6}]_6(r_6)$, that is $[f_6]_6(P_6^*) < r$, $R\left(f_6(P^*)\right) > 1$which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_6(g^{-6}(b))) > 6, i.e.,$

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6\left(g_6^{-6}(b)\right)a_6} > 1. \quad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_6^*(a_6), Q_6^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_6^*(a_6), Q_6^*)$ is more complicated. We again consider perturbations $m_6(t_6 a_6) = p_6(t_6, a_6) - p_6^*(a_6)$ and $W_6(t) = Q_6(t) - Q^*$.Then

$$M_6(t) \equiv \int_0^\infty m_6(t, a_6) da_6 = P_6(t) - P_6^*,$$

$$[m_6]_6 + [m_a]_6 = [p_6]_6 + [p_a]_6 - \left[p^*\right]_6(a_6) \approx -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6$$

$$-[f_{6*}^6]_6 p_6^*(a_6) M - [f_{2*}^6]_6 p_6^*(\mathrm{a}) W_6,$$

$$m_6(t,0)p_6(t,0) - p_6^*(0) \approx h_* \int_0^\infty \beta_6(a_6) m_6(t, a_6)\, da + \frac{[h_*^6]_6}{[h_*]_6}\, p_6^*(0) W_6$$

$$\frac{dW_6}{dt} = \frac{dQ_6}{dt} \approx [g_*^6]_6 Q_6^* M_{,6}$$

Where $[f_{6*}]_6$ denoting $[f_6(P^*)]_6$, $[f_{6*}]_6$ denoting $[f_2(P^*)]_6$, $\left[f_{6*}^{'}\right]_6$ denoting $\left[f^{'}{}_6(P^*)\right]_6 \left[f_{2*}^{'}\right]_6$ denoting $[f_2]_6$ $(Q_6^*,$ $h_{*}^{}{}_6$ denoting $h_*^{'}$ denoting $h_6^{'}$ $(Q_6^*)$, and $\left[g_*^{'}\right]_6$

Denoting $\left[g^{'}\right]_6(P_6^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_6]_6 + [m_a]_6 = -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6 - [f'_{6*}]_6[p_6^*](a_6)M_6 - [f'_{2*}]_6 p_6^*(a_6)W_6$$

$$m_6(t,0) = [h_*]_6 \int_0^\infty \beta_6(a_6)m_6(t,a_6)da_6 + \frac{[h_*^6]_6}{[h^*]_6}[p^*]_6(0)\,W_{,6} \qquad (30)$$

$$\frac{dW_6}{dt} = [g_*^6]_6[Q_6^*]M_6$$

Integrating $[m_6]_6(t,a_6)$ with respect to a from 0 to ∞, we have

$$\frac{dM_6(t)}{dt} = \int_0^\infty [mt]_6(t,a_6)da_6$$

$$= -\int_0^\infty [m_a]_6\,(t,a_6)da_6 - \int_0^\infty \mu(a_6)m(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + [f'_{6*}]_6 \int_0^\infty p_6^*(a_6)da_6\right)M_6 - [f_{2*}^6]_6 \int_0^\infty p_6^*(a_6)da_6 W_6$$

$$= m_6(t,0) - \int_0^\infty \mu_6(a_6)\, m_6(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + [f'_{6*}]_6 P_6^*\right)M_6 - [f'_{2*}]_6 P_6^* W_{,6}$$

From (30), we arrive at the equations

$$\frac{dW_6(t)}{dt} = \int_0^\infty ([h_*]_6\beta_6(a_6) - \mu_6(a_6))m_6(t,a_6)da_6$$

$$-\left([f_{6*}] + [f_{2*}]_6 + [f'_{6*}]_6 P_6^*\right)M_6 + \left(\frac{[h'_*]_6}{[h_*]^6}\,P_6^*(0) - f'_{2*}P_6^*\right)W_6,$$

(36)

$$\frac{dW_6(t)}{dt}\left[g^{'}{}_{*}\right]_6 Q_6^{*} M_{.6}$$

By the mean value theorem,

$$\int_0^{\infty}([h_{*}]_6\beta_6(a_6)-\mu_6(a_6))m_6(t,a_6)da_6=([h_{*}]_6\beta_6(\theta_6)-\mu_6(\theta_6))M_{,6}$$

Where $0<\ \theta_6=\theta_6(t)<\infty$.

And if $a_6^2\,[a_6^2]_6<4\,[a_2]_6<4\,a_2 e$,

$\mathrm{Re}\,\lambda_6\left[\lambda_6\right]_6=\mathrm{Re}\lambda_2\left[\lambda_2\right]_6=-\frac{6}{2}a_6[a2\ ]_6<0.$

Let $\Delta=\max_{2<u<\infty}\left|h_{*}\,[h_{*}]_6\beta_6(a_6)-\mu_6(a_6)\right|$. Then$|n_6(t)|$ 2Δ∀t 0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^{*}(a),\boldsymbol{Q}^{*})$ is asymptotically stable if either $a_6^2<4\,a_2 e$ and $[a_6]>2\Delta$,or $[a_6^2]$ 4$[a_2]$e and $(a_6-\sqrt{(a_6^2)_6-4(a_2)_6 e})>2\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a)=\mu_0+\left(\frac{a}{a+6}\right)$ and $\beta_6(a_6)=ae^{-ma}$ Then $\Delta=[\mu_0]+1$ Simple calculation leads to

$$\int_0^{\infty}\beta_6(a_6)e^{-\int_0^a\mu_6(\tau_6)+d\tau_6}da=\frac{6}{(m+\mu_0+1)^2}+\frac{2}{(m+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_6 + [\mu_0]$ 0.5. Assume$[f_6](x_6)= K_6x_6$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_6(x_6)= [0x_6]/(1 + x_6)$; the functional responses $[[f_2](x) = [0]\, x_6/(1 + x_6)$and $h_6(x_6)$=(1+$\omega_6 x_6$)/(1 + $x_6$)(See Holling [7] and Rashevsky [13])

Equation (64),now has the form of

$$\frac{m_6 + [\mu_0] + [\mathrm{K}][P_0])}{[m_6]_{\,6} + (\mu_0)_{\,6} + [\mathrm{K}][P_0])\,{}^{3}_{6}\ [\mathrm{K}]_{\,6}[P_0]\,{}^{3}_{6}} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(6 + P_0) < b_6$i e $[P_0] < b/(\delta_6 - b_6)$then the equilibrium$(P_0)$ $[P_0]_{\,6}(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_6$

$[f_{6*}] = K_6\ \mathrm{P}^{*}\ [f_{2*}] = \frac{\sigma \mathrm{Q}_6^*}{6+\mathrm{Q}_6^*}$ $[h_*] = \frac{6+\omega_6 \mathrm{Q}_6^*}{6+\mathrm{Q}_6^*}$

$[f_{6*}] = K_6\ \left[\mathrm{f}\,{}^{6}_{2*}\right] = \frac{\sigma_6}{(6+\mathrm{Q}_6^*)^2}[\ \mathrm{g}\,{}^{6}_{*}] = \frac{\sigma_6}{(6+\mathrm{P}_6^*)^2}$ $\left[\mathrm{h}\,{}^{6}_{*}\right] = \frac{\omega_6 - 6}{(6+\mathrm{Q}_6^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0$, K, b, δ, and $\omega$ need to satisfy

$$\frac{\omega_6}{(m + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 + b_6} + \sigma)^{\,2}} + \frac{2\omega_6}{([m_6]_{\,6} + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 + b_6} + \sigma_6)^{\,3}} < 1$$

And

$$\frac{1}{(m_6 + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 - b_6})^{\,2}} + \frac{2\omega_6}{(m_6 + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 - b_6})^{\,3}} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_6 + [\mu_0]_6 + 1 + \frac{K_6 b_6}{\delta_6 - b_6} + (m_6 + [\mu_0]_6 + \frac{K_6 b_6}{\delta_6 - b_6} + \sigma_6) Q_6^*)^3$$

$$= m_6 + [\mu_0]_6 + 3 + \frac{K_6 b_6}{\delta_6 - b_6}(1 + Q^*)(1 + \omega_6 Q_6^*$$

Now

$$\left[a_6^2\right] - 4\left[a_2\right]e = \left(\frac{2K_6 b_6}{\delta_6 - b_6}\right)^2 + \frac{4K_6 b_6 \sigma_6 Q_6^*}{(\delta_6 - b_6)(1 + Q_6^*)} + \frac{\sigma_6^2 Q_6^{*2}}{(1 + Q_6^*)^2}$$

$$\frac{4 b_6 \sigma_6 (\delta_6 - b) Q^*}{\delta_6 (1 + Q_6^*)^2} - \frac{4 (\delta_6 - b_6)^2 (1 - \omega_6) Q_6^*}{\delta_6 (1 + \omega_6 Q_6^*) 1 + Q_6^*)} P * (0)$$

And

$$\Delta = \lim_{0 < a < \infty} |[h_*]_6 a e_6^{mu} [\mu_0]_6 - \frac{a_6}{a_6 + 6}$$

Hence if $[a_6^2 > 4(a_2)_6 e$ and

$$[a_6] = \frac{2K_6 b_6}{\delta_6 - b_6} + \frac{\sigma_6 Q_6^*}{6 + Q_6^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_6]_6 = [f_{6*}]_6 + [f_{2*}]_6 + [f_{6*}]\,P^*_6$$

$$[a_2]_6 = \left[f^6_{2*}\right]P^*_6 - \frac{h^6_*}{h_*}P^*(0)$$

$$e=[g^6_*]_6\,Q^*_6$$

$$\eta_6(t) = \left[h_*\right]_6\beta_6\left(\theta_6(t)\right) - \mu_6\left(\theta(t)\right),$$

The system(36)can be written as the following system:

$$[M(t)]\left[a_6\right]M_6 - [a_2]_6 W_6 - [n(t)]M_6$$

$$w_6(t) = eM_6$$

Write

$$A=\begin{pmatrix} -[a_6]_6 - [a_2]_6 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_6 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_6 \\ W_6 \end{pmatrix}$$

The solution of (32)is

$$X_6(t) = e^{A,t}\,[X]_6(0) + \int_0 e^{At}$$

Other problems to which the model is applied are: To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case β→∞ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness condition of (e) entropy. It is shown that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology

(gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure authors have discovered. It is shown that Legendre invariant thermodynamic scalar curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws, leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions,**

**geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged AdS black holes by considering the case of the RN-AdS black holes surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase transition. For the case $\omega_q=-2/3$ authors derive analytic expressions of (e) critical physical quantities, while (e) for cases $\omega_q\neq-2/3$ appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark

energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=)

chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross

section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions, while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center

collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma? Paul M. Chesler STOP Paul M. Chesler, Krishna Rajagopal** calculate how the energy and (e&eb) the opening angle of jets in N=4 SYM theory evolve as they propagate through (e&eb) the strongly coupled plasma of that theory. **Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb)

heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version) **Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb) strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion

collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar (=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI: 10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills

(SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the

decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed

black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to

equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a

nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about

the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic

map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstates thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11 August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical

advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.**

Characterizing the behaviour of strongly coupled quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at $\lambda=\lambda c$, the correlation length diverges as $\xi\sim|\lambda-\lambda c|-\nu$, while (e) the gap between the ground state and first excited state

vanishes as (e) ΔE∼ξ−z∼|λ−λc|zν. Exactly at the critical point λ=λc, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as ω∼|k−k0|z. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004 STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting

qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element. The basic architecture presented here can be expanded, offering (eb) the possibility for the coherent interaction of (e&eb) a large number of superconducting qubits. **Nature 449, 438-442 (27 September 2007) | doi:10.1038/nature06124; Received 18 April 2007; Accepted 25 July 2007 Coherent quantum state storage and transfer between two phase qubits via a resonant cavity Mika A. Sillanpää1, Jae I. Park1 & Raymond W. Simmonds1 STOP. Chui-Ping Yang, Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e)

long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007**

## PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION ATTRITION MODELS

## MOODULE NUMBER SEVEN

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Systems to which have applied the model are as follows: Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures have been proposed. An attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in**

**(eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP. In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011 http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 Qubit transfer**

between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys. Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer. Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP. In recent years S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P have investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys.

**Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP. Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2 [quant-ph] for this version) Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP. Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue from over (e&eb) a thousand plant and animal samples. NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP. Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this behavior, enhancing and accelerating (eb, eb+) the growth of correlations. Michael Kastner gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. Michael Kastner tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP. A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E.**

**Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91, 075435 – Published 26 February 2015 STOP. Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1 proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point'' with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. Membranes at quantum criticality Petr Hořava1, 2 Published 4**

**March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that (eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035 Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013 STOP Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study**

**the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen STOP Authors study the dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing**

**seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) Dynamic Service Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by H∗(X,(Λ∗TX)⊗g), where X is the target hyperk\"ahler manifold. We further give a mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the**

**low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition. Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb) the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537 © 2016 American Physical Society Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant (Δα/α) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and Δα/α to (e&eb) this conclusion.**

**However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts z≲3, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the Λ cold dark matter model value (jΛCDM=1); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in α(z), linked to (e&eb) a linear coupling (ζ) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order |ζ|∼10−4, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in Δα/α (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems, authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them**

**are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP Authors measure the star formation properties of (e) two large samples of galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the star formation rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a function of environment. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]. Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size √ n × √ n? The problem is that √n time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time O (√n) for d ≥ 3, or O (√nlog5/2n) for d = 2. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ∗This work was mostly done while the author was a PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship and by CIAR. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In**

**d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e) sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP The relationship between the foundations of mathematics and (e&eb) physics is a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb)**

**flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph] (or arXiv: 1503.05600v1 [math-ph] for this version) Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative**

**real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to the fine structure constant. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor goes to 0 as the (e&eb) time goes to 0, the big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph]**

**(or arXiv: 1306.4613v1 [quant-ph] for this version) Effects of gauge theory based number scaling on geometry Paul Benioff STOP In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x, outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph] (or arXiv: 1211.3381v2 [quant-ph] for this version) Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value, \int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances, the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an**

**observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include (e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP This work is based on (e) the field of reference frames based on (e) quantum representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul Benioff STOP A quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of**

**these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference: Journal of Physics: Conference Series 70 Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows: System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The**

**Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol**

## 50. MODEL FORMULATION

Let $p_7(t, a_7)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(g_7)Q_7$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_7(0) = 0, \quad g_7(P_7) > 0, \quad P_7 \text{ 0}, \qquad \lim_{P_7 \to \infty} g_7(P_7) = (K_7)_7 \text{ ∞}. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_7 + (\rho_a)_7 + D_7(a_7, P_7, Q_7)\rho_7 + 0, \qquad 0 < a_7 \text{ ∞},$$

$$\rho_7(t,0) = \int_0^\infty n_7(a_7, P_7, Q_7)\rho_7(t_7, a_7)da_7$$

Where $D_7$the mortality or debit type function is function and $n_7$ is the fertility function.

Put $_7$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_7}{dt}\right]_7 = x_7\left[\frac{dQ_7}{dt}\right]_7 - x_7\left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_7}{dt}\right]_7 = x_3\left[\frac{dQ_3}{dt}\right]_7 - x_4\left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_7}{dt}\right]_7 - x_6\left[\frac{dQ_3}{dt}\right]_7$$

Where $x_7, x_7, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_7$ and $g_7(P_7)Q_7$

Now when you add the term $(b_7Q_7)$ with $g_7(P_7)(Q_7)$ we are infect see is that $(b_7Q_7)$ is credited incrementally part of $g_7(P_7)(Q_7)$ or instant nearly, when subtracted from $\frac{dQ_7}{dt}$ , it means that $g_7(P_7)(Q_7)$ is slowly eaten by $\frac{dQ_7}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

Following equation throws light on important issues in the following cases

**Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA**

(3030) (Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor.

(3031) In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface.

(3032) As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone.

(3033) Such planets are all the "right" distance from (e) their star, and hopefully, also have water.

(3034) But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone.

(3035) For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement.

(3036) But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway.

(3037) But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain.

(3038) It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life.

(3039) Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface.

(3040) The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113

**Abstract**

(3041) The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question.

(3042) Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry.

(3043) This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry.

(3044) Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments.

(3045) Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy.

(3046) The electronic and (e&eb) structural variability among them is extraordinary.

(3047) The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations.

(3048) The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface.

(3049) The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers.

(3050) Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed

mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp

**Water waves travel back in time to retrace their ripples By Jacob Aron (New Scientist)**

(3051) Nudge your bath tub just right and (eb) you might end up reversing time.

(3052) Don't worry – you won't end up wet, naked and without a towel in the Cretaceous period. Emmanuel Fort at the Langevin Institute in Paris, France, and his colleagues have developed what they call a "time mirror", which can make water waves rewind into (e&eb) their past. Humans and rubber duckies are (=) affected.

(3053) To understand how this works, imagine dropping a stone in a pond and watching waves ripple out from a central point in an ever-expanding circle. The equations governing the behaviour of waves say that it should also be possible for a ring of ripples to refocus back into a single point, like a reversed movie of the dropped stone. But we never actually see this happening, because it is vanishingly unlikely that an outside-in ring would form spontaneously. Put it in reverse

(3054) Fort and others have previously shown that it is possible to engineer the situation so the waves turn back on themselves by kitting out the edge of your pond with sensors and wave generators, effectively recording and then playing back the waves in reverse. The process has been done with water, sound and even microwaves, and is used to produce clearer ultrasound images for medical uses.

(3055) But now the team has discovered that simply dropping the tub downwards is enough to send water waves scurrying in reverse, with no need for complex sensors. That's because the velocity of water waves depends on acceleration due to gravity. "If you do a sudden jolt, you change the effective gravity, and the velocity of the wave changes," says Fort.

(3056) The disruption essentially resets the starting point of the wave. Once normal gravity is restored, it starts to travel both backwards and forwards. "Suddenly, you freeze the wave. When you release it again, it doesn't remember that it was going one way," says Fort.

(3057) Ripples in the water returned exactly to their start point - the shape of the Eiffel Tower

(3058) Ripples in the water returned exactly to their start point – the shape of the Eiffel Tower

**Déjà vu**

(3059) To show that the backward-travelling waves converge exactly on their starting point, just as if you were rewinding time, the team conducted a series of experiments. In some, they simply created a circular ripple of waves by dipping a probe in the water, similar to dropping a stone in the pond. In others, they used wave emitters in the shape of France, the Eiffel Tower and a smiley face to create complex wave-fronts. In each case, the reversed waves converged to recreate their original shape.

(3060) The experiments aren't truly reversing time, because the waves continue to travel out again once they have converged at the central point, points out Fort. "A real time-reversal would have a sink that absorbs all the waves," he says.

(3061) Now the team hopes to conduct similar experiments with other kinds of waves, like light or sound, although it may prove more difficult to find an analogue of their time-mirror jolt in other wave mediums, he says.

(3062) "The work is spectacular," says Geoffrey Edelmann of the US Naval Research Laboratory in Washington, D.C. "They essentially change the rules that govern how the wave propagates and, in doing so, create an entirely new way to focus waves free of spatial constraints." He also thinks the technique could be extended to other kinds of waves, and perhaps have medical applications such as lithotripsy, which uses ultrasound waves to break up kidney stones. Journal reference: Nature Physics, DOI: 10.1038/nphys3810

**Time reversal and holography with spacetime transformations Vincent Bacot, Matthieu Labousse, Antonin Eddi, Mathias Fink & Emmanuel Fort Nature Physics (2016) doi: 10.1038/nphys3810**

(3063) Wave control is usually performed by (e) spatially engineering the properties of a medium.

(3064) Because time and space play similar roles (e&eb) in wave propagation, manipulating time boundaries provides (eb) a complementary approach.

(3065) Here, authors experimentally demonstrate the relevance of this concept by (e) introducing instantaneous time mirrors.

(3066) Authors show with water waves that a sudden change of (e) the effective gravity generates (eb) time-reversed waves that refocus at the source.

(3067) They generalize this concept for (e) all kinds of waves, introducing (eb) a universal framework which explains (eb) the effect of any time disruption on (e&eb) wave propagation.

(3068) Authors show that sudden changes of the medium properties generate (eb) instant wave sources that emerge (eb) instantaneously from the entire space at (eb) the time disruption.

(3069) The time-reversed waves originate from (e) these 'Cauchy sources', which are (+) the counterpart of Huygens virtual sources on (eb) a time boundary. It allows us to revisit the holographic method and introduce a new approach for wave control.

**Viewpoint: A Bird's Eye View of Circuit Photons Sebastian Schmidt, Institute for Theoretical Physics, ETH Zurich, 8093 Zurich, Switzerland June 30, 2016• Physics 9, 74 A scanning probe detects the quantum states of photons in a microwave circuit, providing the information needed for quantum simulations. Figure caption expand figure A. Houck/Princeton (deleted)**

(3070) Figure 1: Scanning defect microscopy provides a map of photons in a resonator lattice. Houck and colleagues demonstrated the technique using (e) 49 resonators (grey lines) that were coupled together to (e&eb) form a kagome lattice.

(3071) Quantum mechanics rules the dynamics of (e) light and matter.

(3072) Yet performing a quantum-mechanical simulation of a material from first principles is (=) practically impossible on a classical computer because (e) the complexity of the simulation increases (eb) exponentially with the number of particles involved.

(3073) The solution, according to Richard Feynman, was to build a machine out of (e) quantum building blocks that could directly emulate (e&eb) the material itself [1].

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| (3074) Prototypes of such quantum simulators that are based on (e) ultracold atoms, ions, photons, and superconducting microwave circuits are now available [2], with (e&eb) the latter, in particular, having attracted (e) Silicon Valley's interest. |
| (3075) The challenge with these circuit-based simulators, however, is that they are (=) 2D, which complicates the readout of their constituent elements. Andrew Houck from Princeton University, New Jersey, and colleagues have now delivered an attractive solution by developing a technique [3], called scanning defect microscopy (Fig. 1), that determines (eb) the number of photons occupying (e) each mode of a 2D microwave circuit. |
| (3076) It is this information that would serve as (=) the fundamental input and output for (e) certain quantum simulations. |
| (3077) Superconducting microwave circuits combine electronic and (e&eb) photonic degrees of freedom [4, 5]. |
| (3078) The main element of the circuit is (=) a transmission line, which is made up of (e) a central superconducting wire separated by a gap from (e) two grounded plates. |
| (3079) All of these structures are on a single plane, as if one had taken a 2D slice through (e&eb) a coaxial cable. |
| (3080) When truncated, the transmission line becomes (=) a resonator, which can host (e&eb) discrete photon modes within (eb) its gaps |
| (3081) Large lattices of resonators can be engineered in various 1D or 2D geometries by coupling (e&eb) two, three, or more resonators together via (e&eb) a capacitive interface. In many ways, photons in such devices behave similarly to electrons in a solid. |
| (3082) To make microwave circuits that can simulate (e&eb) quantum phenomena faster than a classical computer, however, resonator lattices have to be integrated with (e&eb) superconducting qubits |
| (3083) . Such qubits are controlled with (e&eb) electrical currents in a Josephson tunnel junction, and in many respects, they behave (e&eb) like artificial atoms, which couple to (e&eb) the photons in the circuit. |
| (3084) As a result, superconducting microwave circuits can be used to (e) explore the coupling between the quantum states of light and (e&eb) matter, the regime of (e) circuit quantum electrodynamics (cQED). |
| (3085) Photons in these devices often exhibit (eb) striking matter-like behavior [6, 7], providing (eb) the basis for the simulation of complex materials. |
| (3086) Such circuits can be fabricated on a substrate using (e) standard lithographic techniques, with (e&eb) qubits and resonators that are (=) hundreds of micrometers or even millimeters in size. |
| (3087) Small-scale realizations of cQED lattices have already been used for (e) proof-of-principle quantum simulations [8, 9]. |
| (3088) However, future large-scale simulations will have to be initialized and measured (eb) efficiently—steps known as (=) drive and readout. |
| (3089) This requires connecting the lattices to (e) external circuitry, such as voltage sources, circulators, and amplifiers, which typically reside outside (e) the chip. |
| (3090) So far, this connectivity has been realized through (e&eb) external transmission lines, but these elements can only couple to (e&eb) the resonators at the edges of (e) the lattice because of (e) its 2D nature. |
| (3091) Such a side view provides (eb) only limited information about (e&eb) the quantum state and spatial correlations of photons inside (eb) the bulk resonators, seriously limiting (e) the applicability of circuit QED devices for (e) quantum simulations. |

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| (3092) The probe developed by Houck and colleagues, scanning defect microscopy (SDM), solves (e) this problem by (e) moving outside the plane of the chip to provide (eb) a fuller image of photons within (eb) the chip. |
| (3093) In SDM, a small sapphire tip is positioned a few tens of (e) micrometers above the surface of one of the resonators (Fig. 1). |
| (3094) The presence of this dielectric modifies (e&eb) the evanescent electric field from the circuit photons, which extends (eb) just outside the plane of the chip. |
| (3095) Specifically, the tip increases (eb) the local capacitance of the transmission line, downshifting (e) the local resonator frequency. |
| (3096) This in turn causes (eb) a drop in the transmission of photons through (e&eb) the lattice when (e) it is driven at its nominal resonance frequency. |
| (3097) Houck's team showed that if the sapphire tip has only a weak (linear) effect on the resonator frequency, then (eb) a simple analytic relationship exists between the induced frequency shift, (e&eb) the change in transmission, and (e&eb) the local photonic occupation of each resonator. |
| (3098) This relationship provides (eb) a recipe for detecting the local photon occupancy inside (eb) the lattice. |
| (3099) First, calibrate the probe by determining (eb) the frequency shift for (e) a single isolated resonator as (=) a function of the tip's position. |
| (3100) Second, measure the transmission through (e&eb) the whole lattice with and without (e&eb) the probe. |
| (3101) Finally, use the measured transmission drop with (e&eb) the probe present to find (eb) the photonic occupation of (e) an individual resonator in (eb) the lattice. |
| (3102) A map of the photon-mode occupancies can be obtained by (e) scanning the probe across the whole lattice. |
| (3103) Houck's team used (e) SDM to image the occupancies of (e) the photon modes of (e) 49 coupled resonators arranged in (eb) a special geometry called a kagome lattice (see details in Fig. 1). |
| (3104) (To simplify the measurement, no superconducting qubits were (=) incorporated with (e&eb) the resonators.) |
| (3105) The occupancies they measured were in excellent agreement with (=) theoretical predictions, corresponding to (e&eb) a fidelity of 0.99. |
| (3106) The high fidelity is even more impressive considering the weak signal-to-noise ratio of (e) their technique (which is limited by (e) the intentionally small size of the frequency-shifted region) and (e&eb) the technical challenge of (e) positioning the probe inside (eb) a cryostat for (e) in situ measurements. |
| (3107) While SDM shares similarities with (e&eb) so-called bead-pull experiments for accelerators and scanning gate microscopy for (e) nanostructures, it is truly a (=) new tool for quantum electronics. |
| (3108) First application of SDM is likely to be (=) the simulation of 2D topological phenomena and (e&eb) geometrically frustrated lattices (of which the kagome lattice is (=) one example.) |
| (3109) Both are associated with (e&eb) macroscopically degenerate systems, which are (=) notoriously difficult to simulate on (e&eb) a classical computer. |
| (3110) Such circuits are (=) genuinely open systems (photons can be added or lost). |
| (3111) This combined with (e&eb) their inherent strong light-matter coupling, makes (eb) such quantum simulators ideal candidates for (e) the study of many-body, out-of-equilibrium phenomena, such as (=) thermalization or nonequilibrium phase transitions. |

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| (3112) SDM could also lead to (eb) the discovery of new exotic phenomena, which are (=) unknown in traditional solid-state systems. |
| (3113) In a nutshell, SDM constitutes a new and indispensible tool for (e) future experiments on quantum simulations with microwave photons. This research is published in Physical Review X. |

**Imaging Photon Lattice States by Scanning Defect Microscopy D. L. Underwood, W. E. Shanks, Andy C. Y. Li, Lamia Ateshian, Jens Koch, and A. A. Houck Phys. Rev. X 6, 021044 – Published 30 June 2016**

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| (3114) Microwave photons inside (eb) lattices of coupled resonators and superconducting qubits can exhibit (eb) surprising matter like behavior. |
| (3115) Realizing such open-system quantum simulators presents (eb) an experimental challenge and requires (e) new tools and measurement techniques. |
| (3116) Here, authors introduce scanning defect microscopy as one such tool and illustrate its use in mapping (e&eb) the normal-mode structure of microwave photons inside (eb) a 49-site kagome lattice of coplanar waveguide resonators |
| (3117) Scanning is accomplished by (e) moving a probe equipped with a sapphire tip across (e&eb) the lattice. |
| (3118) This locally perturbs (e&eb) resonator frequencies and induces (eb) shifts of the lattice resonance frequencies, which authors determine by (e) measuring the transmission spectrum. |
| (3119) From the magnitude of mode shifts, author can reconstruct (eb) photon field amplitudes at (eb) each lattice site and thus create (eb) spatial images of the photon-lattice normal modes. Received 4 December 2015 DOI:http://dx.doi.org/10.1103/PhysRevX.6.021044 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society |

| |
|---|
| |
| We assume that |
| $$D_7 = \mu_7(a_7) + (f_7)_7(P_7)_7 + (f_7)_7(Q)_7 \quad (7)$$ |
| $$n_7 = \beta_7(a_7)h_7(Q_7), \quad (3)$$ |
| Where $\mu_7$the age- dependent debit rate whit is $\mu_7(a_7)$ ▯ $\forall a > 0$; $(f_7)_7$ is the density- dependent death rate with $f_7(0) = 0$, $(f_7^7)_7(P_7)0$, $\forall P$ ▯0, and $\lim_{p\to\infty}(f_7)_7(P_7) = (k_7)_7$ ▯∞; $f_7$ is the functional response of creditors to prey with $[f_7]_7(0) = 0, (f_7^7)_7(Q) > 0, \forall Q$ ▯, and $\lim_{Q\to\infty}(f_7)_7(Q_7) = [k_7]$ ▯∞; $\beta$ is the age dependent birth rate with $\beta(a)$ ▯$, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^7(Q)$ ▯$, \forall Q$ ▯0, and $\lim_{Q\to\infty} h(Q) = k_3$ ▯. |

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_7 + (\rho_a)_7 = -\mu_7(a_7) + (f_7)_7 P_7) + \big(f_7)_7(Q_7)\big)(\rho_7)_7 \quad (4)$$

$$\rho_7(t_7, 0) = h_7(Q_7) \int_0^\infty \beta_7(a_7)\rho_7(t, a_7)da, \quad (5)$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(P_7)Q_7 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_7$ and $Q_7$ of (4-6) satisfy

$$0 \leq \rho_7(t, a_7) \leq \rho_7^*(t, a_7),$$

$$0 \leq Q_7(t) \leq Q_7^*(t),$$

Where $\rho_7^*(t, a_7)$ and $Q_7^*(t)$, are solutions of

$$(\rho_t^*)_7 + (\rho_a^*)_7 = -(\mu_7(a_7) + (f_7)_7(P_7^*) + (f_7)_7 Q_7^*))\,\rho_7^*, \quad (7)$$

$$\rho_7^*(t, 0) = \int_0^\infty \beta_7(a_7)\,\rho_7^*(t, a_7)da_7 \quad (8)$$

$$\frac{dQ_7^*}{dt} = -b_7 Q_7^* + g_7(P_7^*)Q_7^*, \quad (9)$$

With $P_7^*(t) = \int_0^\infty \rho_7^*(t, a_7)da_7$

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From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_7^*(t)}{dt} = P_7^*(t)\big(r_7 - (f_7)(P_7^*) - (f_7)_7(Q^*)\big), \quad (10)$$

$$\frac{dQ_7^*(t)}{dt} = Q_7^*(t)(-b_7 + g_7(P_7^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_7(a)e^{-\int_0^a \mu_7(\tau)dt}e^{-ra}\,da_7 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_7^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(t)dt}\,da_7 > 1, \quad (H1)$$

Which guarantees $r_7 > 0$.

*Remark.* If $\beta_7(a) = [\beta_0]_7 a_7 e^{-\alpha_7 a_7}, \mu_7(a_7) \equiv \mu_7$ then the solution of (12) is

$$r_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_7]_7(P_7) = 0, [f_7](Q_7) = \lambda_7 Q_7\ h_7(Q_7) = 1, \quad and\ \ g_7(P_7) = c_7 P_7$$

$$[\rho_7]_7 + [\rho_a]_7 = -(\mu_7 + \lambda_7 Q_7)\rho_7$$

$$\rho_7(t,0) = \int_0^{\infty} [\beta_0]_7 a_7 e^{-\alpha_7 a_7} \rho_7(t,a_7)\, da_7$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + c_7 P_7 Q_7$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_7}{dt} = \omega_7 P_7 - \lambda_7 P_7 Q_7$$

$$\frac{dQ_7}{dt} = b_7 Q_7 - c_7 P_7 Q_7$$

Where $\omega_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$

## 51. THE EQUILIBRIUM $(\rho_0)_7(a_7), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_7(a_7), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_7 + (\rho_a)_7 = -\big(\mu_7(a_7) + (f_7)_7(P_7)\big)\rho,$$

$$\rho_7(t,0)$$

$$= \int_0^\infty \beta_7\,(a_7)\rho_7(t,a_7)da_7 \qquad (13)$$

The solution of (13) has the form of

$$\rho_7(t,a_7) = Ce_7^{r(t-a_7)} - \int_0^a (\tau)d\tau - \int_0^7 (f_7)_7\big(P_7(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_7(t) = C_7 e^{r_7 t - \int_0^t (f_7)_7(P_7(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_7^{-7}(r)$ if $k_7 > r$. This is equivalent to

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{(k_7)_7 a_7} da_7 < 1.$$

In fact, if we define

$$R_7(s_7) = \int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau)d\tau} e^{-s_7 a_7}\, da_7$$

$R(s)$ is a decreasing function. $R_7(r_7) = 1$ And $R_7(k_7)_7 < 1$ imply $r_7 < (k_7)_7$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_7(a_7), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{k_7 a_7} da_7$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_7 \equiv \int_0^\infty (\rho_0)_{77}(a_7)da_7$ *satisfies*

$$\int_0^{\infty} \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{-(f_7)_7(P_0)a_7} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_7(a_7), 0)$

Let $u_7(t, a_7) = \rho(t, a_7) - \rho_0(a_7)$ Then

$$N_7(t) \equiv \int_0^{\infty} u_7(t, a_7) da_7 = \int_0^{\infty} \rho_7(t, a_7) da_7 - \int_0^{\infty} (\rho_0)_7(a_7) da = P_7(t) - P_0)_7$$

And

$$(u_7)_7 + (u_a)_7 = (\rho_t)_7 + (\rho_a)_7 - (\rho_0^7)_7(a_7)$$

$$= -\mu_7(a_7) + (f_7)_7(P_7) + (f_7)_7(Q_7))(u_7 + (\rho_0)_7) + \mu_7(a_7)(\rho_0)_7 + (f_7)_7(P_0)_7(\rho_0)_7$$

$$= -\mu_7(a_7)u_7 - (f_7)_7(P_7) + (f_7)_7(Q_7))\, u_7 - (\rho_0)_7\, (f_7)_7(P_7) - (f_7)_7(P_0)_7$$

$$+ (f_7)_7(Q_7)), \qquad (15)$$

Where $(\rho_0^7)_7(a_7)$ denotes $d(\rho_0)_7(a_7)/da_7$

Expansion $(f_7)_7(P_7)$ and $(f_7)_7(Q_7)$ at $(P_0)_7$ and 0 yield respectively

$$(f_7)_7(P_7) = (f_7)_7((P_0)_7 + N_7)$$

$$= (f_7)_7(P_0)_7 + (f_7^7)_7(P_0)_7 N_7 + o_7(N_7), \qquad (16)$$

$$(f_7)_7(Q_7)$$

$$= (f_7^7)_7(0)Q + o_7(Q_7) \qquad (17)$$

Write $f_7(P_0) = f_{70}, f_7'(P_0) = f_{70}'$, and $f_7'(0) = f_{70}'$ and substitute (16) and (17) into (15) to obtain

$$(u_7)_7 + (u_a)_7 \approx -(\mu_7(a_7) + (f_{70})_7)u - (f_{70}^7)_7(\rho_0)_7(N_7) - (f_{70}^7)_7(\rho_0)_7(Q_7)$$

$$u_7(t, 0) = \rho_7(t, 0)(\rho_0)_7(0)$$

$$= h_7(Q_7)\int_0^{\infty} \beta_7(a_7)(u_7(t,a) + (\rho_0)_7(a_7)da_7 - (\rho_0)_7\,(0)$$

$$= h_7(Q_7)\int_0^{\infty} \beta_7(a_7)(u_7(t,a_7)da_7 + (h_7(Q_7) - 1)(\rho_0)_7\,(0) \qquad (18)$$

Expanding

$$h_7(Q_7) = 1 + \ (h^7)_7(0)Q_7 + 0(Q_7) \equiv 1 + h_0^{'}Q_7 + 0Q_7$$

And substituting it into (18), we arrive at

$$u_7(t,0) \approx \int_0^{\infty} \beta_7\,(a_7)u_7(t,a_7)da_7 + (h_{a_7}^{'})(\rho_0)_7Q_7$$





The linearization is the following :

$$[u_7] + [u_7] = -(\mu_7(a_7) + [t_{70}]_7u_7 - [\alpha_7]_7(a_7)N_7 - [\alpha_7]_7(a_7)Q_7$$

$$u_7(t,0) = \int_0^{\infty} \beta_7(a_7)u_7(t,a_7)da_7 + [\xi_7]Q_7 \qquad (19)$$

$$\frac{dQ_7}{dt} = -b_7Q_7$$

where

$$[\alpha_7]_7(a_7) = [t_7^7a]_7[p,o]_7(a_7)d2(a_7) = [f_{70}^7]_7(\rho\ o\ )_7(a_7)\xi = [\mathrm{h\ o}^7]_7[\rho\mathrm{o}]_7(\mathrm{o}) \text{ and } \mathrm{b}_7 = \mathrm{b}_7 - g_7(P_7)$$

If $g_7(P_0) > b$, clearly, $(\rho_0)_7(a_7), 0)$ is unstable and the predator population is persistent. Assume $g_7(P_0)_7 < b$.

Then, $Q_7(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_7 + u_a = -(\mu_7(a_7) + (f_{70})_7)u_7 - (\alpha_7)_7(a_7)N_7$$

$$= -(\mu_7(a_7) + f_{70})u_7 - \alpha_7(a_7)\int_0^\infty u_7(t, a_7)da_7$$

$$u_7(t, 0) = \int_0^\infty \beta_7\,(a_7)u_7(t, a_7)da_7$$

Using the method developed by Cushing [2, 3], we suppose

$$u_7(t, a_7) = \tilde{u}_7(a_7)^{e^{\gamma(t-a_7)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_7}\,da_7$. Then,

A solution of the first differential equation for $\tilde{u}(a_7)$ gives,

$$\tilde{u}_7(a_7) = e^{-\int_0^a \mu_7(\tau_7)d\tau - (f_{70})_7 a}\left(\tilde{u}_7(0) - w\int_0^a e^{\int_0^\tau \mu_7(s_7)ds + (f_{70})}\tau_7\,(\alpha_7)_7(\tau_7)e^{\gamma_7\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_7(a_7)e^{-\gamma_7 a_7 - (f_{70})_7 a_7 - \int_0^a \mu_7(\tau_7)d\tau}\right)\tilde{u}_7(0)$$

$$+w\int_0^\infty \beta_7(a_7)e^{-\gamma_7 a_7 - (f_{70})_7 a_7 - \int_0^a \mu_7(\tau_7)d\tau}\int_0^a (\alpha_7)_7(\tau_7)e^{-\gamma_7\tau_7}e^{(f_{70})_7\tau_7 + \int_0^\tau \mu_7(s_7)ds}d\tau_7\,da_7$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty} e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau} da\right)(\tilde{u}_7(0))_7$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau}\int_0^a e^{(f_{70})_7\tau_7+\int_0^\tau \mu_7(s_7)ds+y\tau}\alpha_7(\tau)d\tau\, da\right)w=0. \quad (23)$$

Define $Y(\gamma,a)=e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(t_7)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)da_7\right)u_7(0)$$

$$+w\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\, da_7=0, \quad (24)$$

$$\left(\int_0^{\infty}Y_7(\gamma_7,a_7)da_7\right)\tilde{u}_7(0)$$

$$+\left(-1-\int_0^{\infty}Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\, da_7\right)w=0.$$

Equation (20) has a solution of the form of $u_7(t,a_7)=u_7(a_7)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_7(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_7(\gamma_7,a_7)=\int_0^a(\rho_0)_7(\tau_7)Y^{-7}(\gamma_7,a_7)d\tau_7$. We obtain the characteristic equation

$$f_{70_7}^{l}\left(\int_0^{\infty} Y_7(\gamma_7,a_7)da_7\right)\left(\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$+\left(1-\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)da\right)\left(1+(f_{70}^{l})_7\int_0^{\infty}Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_7(a_7)=(\rho_0)_7(0)e^{-\int_0^a \mu_7(\tau_7)d\tau_7-(f_{70})_7a_7}$$

$$I_7(\gamma_7,a_7)=\int_0^a(\rho_0)_7(\tau_7)e^{\gamma_7\tau_7+(f_{70})_7\tau+\int_0^{\tau}\mu_7(s_7)ds_7}d\tau$$

$$=\frac{(\rho_0)_7(0)}{\gamma_7}(e^{\gamma_7a_7}-1)$$

Let

$$E_7(a_7)=e^{-(f_{70})_7a_7-\int_0^a\mu_7(\tau_7)d\tau_7}$$

Then $[Y_7]_7[\gamma_7]_7,a_7)=E_7(a_7)e^{-\gamma_7a_7}$. by a simple calculation,

$$\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da=\frac{[\rho_0]_7(0)}{\gamma}\left(1-\int_0^{\infty}\beta_7(a_7)E_7(a_7)e^{-\gamma_7a_7}\ da\right);$$

$$\int_0^{\infty}Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da=\frac{[\rho_0]_7(0)}{\gamma_7}\left(\int_0^{\infty}E_7(a_7)e^{-\gamma_7a_7}\ da_7\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_7(0)}{\gamma_7}[f_{70}^7]\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\,(\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7)$$

$$+\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right).$$

$$\left(1+\frac{(\rho_0)_7(0)}{\gamma_7}(f_{70}^{'})_7\left(\int_0^\infty E_7(a_7)da_7-\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right)(1+\frac{(f_{70}^{'})_7}{\gamma_7}(\rho_0)_7(0)\int_0^\infty E_7(a_7)da_7)=0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)da_7 = 1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7 \neq 0,$$

Since $\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)\int_0^a \alpha_7(\tau_7)Y^{-7}(\gamma_7,\tau_7)d\tau_7 da_7 > 0$. Therefore, the characteristic solution is

$$\gamma_7 = -(f_{70}^{'})_7(\rho_0)_7\int_0^\infty E_7(a_7)da_7.$$

Since $f_7'(P_0)_7 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_7), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE OF DEBIT AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_7), 0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_7 + (\mu_7(a_7) + [f_7]_7(P_7^*) + [f_7]_7(Q_7^*))P_7^* = 0, \quad (76)$$

$$p_7^*(0) = h_7(Q_7^*)\int_0^\infty \beta_7(a_7)p_7^*(a_J)p_7^*(a_7)da_7 \quad (77)$$

$$g_7(P_7^*) - b_7 = 0, \quad (78)$$

Where $P_7^* = \int_0^\infty p_7^*(a_7)\, da_7$

Equation (28) yields $P^* = g^{-7}(b)$,

From (26),

$$p_7^*(a) = p_7^*(0)e^{\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(Q_7^*)a_7}$$

Substituting into (27), we need to solve

$$\left(1 - h_7(Q_7^*)\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(Q_7^*)a_7}da_7\right)p_7^*(0) = 0$$

For $Q_7^*$

Fix $p_7^*(0)$ and $P_7^*$ Let

$$H_7(z_7) = 1 - h_7(z_7)\int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(z_7)a_7}da_{,7}$$

And assume that

$$[k_3]_7\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(g_7^{-7}(b_7))a_7 - [k_7]_7 a_7}\, da_7 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_7(z_7)] > 0.$$

Clearly,

$$\frac{dH_7}{dz_7} = h_7(z)(\,[f_7']_7(z_7)\,)\int_0^\infty a_7\beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - f_7(P_7^*)a_7 - f_7(z_7)a_7}da_{7}$$

$-h'(z_7)\int_0^\infty \beta_7(a_7)\, e^{\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(z_7)a_7}\quad da_7$

(See author's It is a debit-credit world)

By the assumptions on $\left[f_7^7\right]_7 and\ [h^7]^7 dH_7/dz > 0$. Hence, if

$$H_7(0) = 1 - \int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7} da_7 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_7(0) < 0$ follows if $g_7^-(b_7) < [f_7^{-7}]_7(r_7)$, where $r_7$ is the solution of (12).

Indeed, since

$$R_7(s_7) = \int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 e^{-y\frac{a}{7}}} da_7$$

Is a decreasing function of $s_7$ when $P_7^* = [g^{-7}]_7(b_7) < [f_7^{-7}]_7(r_7)$, that is $[f_7]_7(P_7^*) < r$, $R\left(f_7(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

**THEOREM 3**

Assume that the hypothesis (H3) holds and $R(f_7(g^{-7}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7\left(g_7^{-7}(b)\right)a_7} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_7^*(a_7), Q_7^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_7^*(a_7), Q_7^*)$ is more complicated. We again consider perturbations $m_7(t_7 a_7) = p_7(t_7, a_7) - p_7^*(a_7)$ and $W_7(t) = Q_7(t) - Q^*$.Then

$M_7(t) \equiv \int_0^\infty m_7(t,a_7)da_7 = P_7(t) - P_7^*,$

$$[m_7]_7 + [m_a]_7 = [p_7]_7 + [p_a]_7 - [p^*]_7(a_7) \approx -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7$$

$$-[f_{7*}^7]_7 p_7^*(a_7)M - [f_{7*}^7]_7 p_7^*(\mathrm{a})W_7,$$

$$m_7(t,0)p_7(t,0) - p_7^*(0) \approx h_* \int_0^\infty \beta_7(a_7)m_7(t,a_7)\,da + \frac{[h_*^7]_7}{[h_*]_7}\, p_7^*(0)W_7$$

$$\frac{dW_7}{dt} = \frac{dQ_7}{dt} \approx [g_*^7]_7 Q_7^* M,_7$$

Where $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}']_7$ denoting $[f'_7(P^*)]_7[f_{7*}']_7$ denoting

$[f_7]_7$ $(Q_7^*,$ $h_{*}^{7}$ denoting $h_*'$ denoting $h_7'$ $(Q_7^*)$, and $[g_*']_7$

Denoting $[g']_7(P_7^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$[m_7]_7 + [m_a]_7 = -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7 - [f'_{7*}]_7[p_7^*](a_7)M_7 - [f_{7*}']_7 p_7^*(a_7)W_7$

$m_7(t,0) = [h_*]_7 \int_0^\infty \beta_7(a_7)m_7(t,a_7)da_7 + \frac{[h_*^7]_7}{[h^*]_7}[p^*]_7(0)\,W,_7$ (30)

$$\frac{dW_7}{dt} = [g_*^7]_7[Q_7^*]M_7$$

Integrating $[m_7]_7(t,a_7)$ with respect to a from 0 to ∞, we have

$$\frac{dM_7(t)}{dt} = \int_0^{\infty} [mt]_7(t, a_7)da_7$$

$$= -\int_0^{\infty} [m_a]_7\,(t, a_7)da_7 - \int_0^{\infty}\ \mu(a_7)m(t, a_7)da_7$$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f^{'}{}_{7*}]_7 \int_0^{\infty} p_7^*(a_7)da_7\right)M_7 - [f_{7*}^7]_7 \int_0^{\infty} p_7^*(a_7)da_7 W_7$$

$$= m_7(t, 0) - \int_0^{\infty} \mu_7(a_7)\, m_7(t, a_7)da_7$$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 - [f_{7*}^{'}]_7 P_7^* W_{.7}$$

From (30), we arrive at the equations

$$\frac{dW_7(t)}{dt} = \int_0^{\infty} ([h_*]_7\beta_7(a_7) - \mu_7(a_7))m_7(t, a_7)da_7$$

$$-\left([f_{7*}] + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 + \left(\frac{[h_*^{'}]_7}{[h_*]^7}\ P_7^*(0) - f_{7*}^{'} P_7^*\right)W_7,$$

(31)

$$\frac{dW_7(t)}{dt}[g^{'}{}_*]_7\, Q_7^*\, M_{.7}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{"*}]_7\beta_7(a_7) - \mu_7(a_7))m_7(t, a_7)da_7 = ([h_*]_7\beta_7(\theta_7) - \mu_7(\theta_7))M,_7$$

Where $0< \theta_7 = \theta_7(t) < \infty$.

And if a $_7^7$ [a $_7^7$]2<4 [a $_7$]2<4 a $_7$e,

Re $\lambda_7[\lambda_7]_7 = \mathrm{Re}\lambda_7[\lambda_7]_7 = -\frac{7}{7}a_7[\mathrm{a}2\ ]_7 < 0$.

Let $\Delta$=max$_{7<u<\infty}|\mathrm{h}_*[\mathrm{h}_*]_7\beta_7(a_7) - \mu_7(a_7)|$. Then$|n_7(\mathrm{t})|$ ⍰⊿∀t ⍰. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium($\rho^*$(a), $\boldsymbol{Q}^*$) is asymptotically stable if either a $_7^7$ <4 a $_7$ e and [a $_7$] > 1⊿,or [a $_7^7$] ⍰[a $_7$]e and (a $_7 - \sqrt{(\mathrm{a}_7^7)_7 - 4(\mathrm{a}_7)_7\mathrm{e}} > 1\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{\mathrm{a}}{\mathrm{a}+7}\right)$ and $\beta_7(a_7) = \mathrm{ae}^{-\mathrm{ma}}$ Then ⊿=$[\mu_0] + 2$ Simple calculation leads to

$$\int_0^\infty \beta_7(a_7)\mathrm{e}^{-\int_0^{\mathrm{a}}\mu_7(\tau_7)+d\tau_7\mathrm{da}} = \frac{1}{(\mathrm{m}+\mu_0+1)^7} + \frac{1}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_7 + [\mu_0]$ ⍰.5. Assume$[f_7](x_7)$= $K_7x_7$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_7(x_7)$= $[0x_7]/(2 + x_7)$; the functional responses $[[f_7](x) = [0]\ x_7/(1 + x_7)$and $h_7(x_7)$=$(1+\omega_7x_7)/(1 + x_7)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_7 + [\mu_0] + [\mathrm{K}][P_0])}{[m_7]_7 + (\mu_0)_7 + [\mathrm{K}][P_0])_7^3\ [\mathrm{K}]_7[P_0]_7^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1+P_0) < b_7$ i e $[P_0] < b/(\delta_7 - b_7)$then the equilibrium$(P_0)$ $[P_0]_7(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_7$

$[f_{7*}] = K_7\ \mathrm{P}^*\ [f_{7*}] = \frac{\sigma \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*} \quad [h_*] = \frac{7+\omega_7 \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*}$

$[f_{7*}] = K_7\ [\mathrm{f}_{7*}^{\ 7}] = \frac{\sigma_7}{(7+\mathrm{Q}_7^*)^7} [\ \mathrm{g}_*^{\ 7}] = \frac{\sigma_7}{(7+\mathrm{P}_7^*)^7} \quad [\mathrm{h}_*^{\ 7}] = \frac{\omega_7 - 7}{(7+\mathrm{Q}_7^*)^7}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_7}{(m + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma)^7} + \frac{1\omega_7}{([m_7]_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma_7)^3} < 1$$

And

$$\frac{1}{(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^7} + \frac{1\omega_7}{(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $\mathrm{Q}^*$ which can be solved from the following equation:

$(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7} + (m_7 + [\mu_0]_7 + \frac{K_7 b_7}{\delta_7 - b_7} + \sigma_7)\ \mathrm{Q}_7^*)^3$

$= m_7 + [\mu_0]_7 + 3 + \frac{K_7 b_7}{\delta_7 - b_7}(1 + \mathrm{Q}^*)(1 + \omega_7 \mathrm{Q}_7^*$

Now

$$\left[\mathrm{a}\begin{matrix}1\\1\end{matrix}\right]-4\left[\mathrm{a}_{7}\right]\mathrm{e}=\left(\frac{1K_7 b_7}{\delta_7-b_7}\right)^{7}+\frac{4K_7 b_7\sigma_7\,\mathrm{Q}_{7}^{*}}{(\delta_7-b_7)(7+\mathrm{Q}_{7}^{*})}+\frac{\sigma_{7}^{7}\mathrm{Q}_{7}^{*7}}{\left(1+\mathrm{Q}_{7}^{*}\right)^{7}}$$

$$\frac{4\,b_7\sigma_7\,(\delta_{7}-\mathrm{b})\mathrm{Q}^{*}}{\delta_{7}\,(1+\mathrm{Q}_{7}^{*})^{7}}-\frac{4\,(\delta_{7}-b_7)^{7}\,(1-\omega_7)\mathrm{Q}_{7}^{*}}{\delta_{7}\left(1+\omega_7\mathrm{Q}_{7}^{*}\right)(1+\mathrm{Q}_{7}^{*})}\mathrm{P}*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_{7}\,ae_{7}^{mu}\;\;[\mu_0]_{7}-\frac{a_7}{a_7+7}$$

Hence if $[\mathrm{a}_{7}^{7}>4(a_7)_{7}\mathrm{e}$ and

In the following are given some more modular structures that are applied with the model and the concomitant results derived thereof:

Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments

(1) We have an intuitive understanding of how the world works. Our intuition works well in our everyday lives, but it doesn't quite match (e&eb) reality.

(2) One of the basic assumptions authors make about objects is that they are (=) objectively real.

(3) For example, when you leave your apartment in the morning, you imagine your house is right where you left it as you go about your day. Barring some calamity, you expect to find your apartment just as you left it when (e) you return in the evening.

(4) This idea is known as Macrorealism, in that a macroscopic object such as your house is still there even when (e) you're not looking at it

(5) In scientific terms we would say the state of an object exists independent of (e) an observer. While this seems obviously true, we can test the idea experimentally.

(6) To show this, let's consider an experiment. Suppose we have a box that contains a coin. At any time we can open the box and (eb) see whether it's heads or tails.

(7) The catch is (=) that the box has to follow a certain rule. Whatever you observe initially, (eb) you must observe the same thing ten seconds later.

(8) If you open the box and see heads, then (eb) ten seconds later when you open the box you must also see heads.

$$[a_7] = \frac{1K_7 b_7}{\delta_7 - b_7} + \frac{\sigma_7 \, \mathrm{Q}_7^*}{1 + \mathrm{Q}_7^*} > 1\vartriangle$$

This equation is important in its consequence in respect of following cases for which the model is applied:

**Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments**

(31) For the experiment we'll open the box, then open it five seconds later, then finally open it again at ten seconds. We know the first and last observations have to agree, but what about the observation in the middle? It depends upon what the box does with (e&eb) the coin.

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| (32) One possibility is that the box does nothing, so if you see tails the first time (eb), you will see tails both other times. |
| (33) Another possibility is that each time you open the box (eb) it flips the coin when you close it. |
| (34) So the first time you see heads, (eb) the second time tails, and (eb) the third time heads again. |
| (35) But they could also make a tricky box that flips (e&eb) the coin at seemingly random times, but does so in such a way that every ten seconds it repeats (eb) the pattern. In that case, authors would always see the same thing the first and last time, but the middle observation would seem (eb) random. |
| (36) If Macrorealism is true, then (eb) middle observations being always the same as (=) the first and last observation, always opposite or always random are (=) the only possibilities for the experiment |

**Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry Download as PDF (Size: 62KB) HTML PP. 50-53 DOI: 10.4236/jqis.2011.12007 5,508 Downloads 14,712 Views Citations Author(s) Leave a comment Mohamed S. El Naschie**

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| (31) Building upon the pioneering work of J. Bell [1] and an incredible result due to L. Hardy [2] it was shown (eb) that the probability of quantum entanglement of two particles is (=) a maximum of 9.0169945 percent [2]. |
| (32) This happens to be exactly the golden mean to (e) the power of five (?5) [3-7]. |
| (33) Although it has gone largely unnoticed for a long time, this result was essentially established independently in a much wider context by (e) the present author almost two decades ago [3-6]. |
| (34) The present work gives two fundamentally different derivations of (e) Hardy's beautiful result leading to (eb) precisely the same general conclusion, namely that by virtue of (e) the zero measure of the underlying |
| (35) Cantorian-fractal spacetime geometry the notion of (e) spatial separability in quantum physics is (=) devoid of any meaning [7]. The first derivation is purely logical and uses (e) a probability theory which combines the discrete with (e&eb) the continuum. |
| (36) The second derivation is purely geometrical and topological using (e) the fundamental equations of a theory developed by the author and his collaborators frequently referred to as E-infinity or Cantorian spacetime theory [3-7]. KEYWORDS Hardy's Quantum Entanglement, Golden mean, Cantor sets, Fractal Spacetime, E-Infinity Theory, Quantum Mechanics, J. S. Bell Cite this paper M. El Naschie, "Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry," Journal of Quantum Information Science, Vol. 1 No. 2, 2011, pp. 50-53. Doi: 10.4236/jqis.2011.12007. |

**Chaos, Solitons & Fractals Volume 28, Issue 4, May 2006, Pages 857–878 Sir Hermann Bondi 1919–2005 Cantorian spacetime and Hilbert space: Part I—Foundations G. Iovane http://dx.doi.org/10.1016/j.chaos.2005.08.074**

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| (51) Authors are going to show the link between the ε(∞) Cantorian space and( e&eb) the Hilbert spaces H(∞). |
| (52) In particular, El Naschie's ε (∞) is (=) a physical spacetime, i.e. an infinite dimensional fractal space, where (e) time is spacialized and the transfinite nature manifests (eb) itself. |
| (53) El Naschie's Cantorian spacetime is an arena where (e) the physics laws appear at each scale in a self-similar way linked to (e&eb) the resolution of the act of observation. |
| (54) By contrast the Hilbert space H (∞) is (=) a mathematical support, which describes (eb) the interaction between the observer and (e&eb) the dynamical system under (e&eb) measurement. |

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| (55) The present formulation, which is based on (e) the non-classical Cantorian geometry and topology of spacetime, automatically solves (e) the paradoxical outcome of (e) the two-slit experiment and the so-called particle–wave duality. |
| (56) In particular, measurement (i.e. the observation) is equivalent to (=) a projection of ε(∞) in the Hilbert space built on 3 + 1 Euclidean spacetime. |
| (57) Consequently, the wave–particle duality becomes (=) a mere natural consequence of (e) conducting an experiment in (eb) a spacetime with (e&eb) non-classical topological and (e&eb) geometrical structures, while (e) observing and taking measurements in (eb) a classical smooth 3 + 1 Euclidean spacetime. |
| (58) In other words, the experimental fact that a wave–particle duality exists is (=) an indirect confirmation of the existence of ε(∞) and a property of (e) the quantum–classical interface. |
| (59) Another direct consequence of the fact that real spacetime is) =) infinite dimensional hierarchical ε(∞) is (=) the existence of scaling law R(N), introduced by the author, which generalizes the Compton wavelength. |
| (60) It gives an answer to the problem of (e) segregation of matter at different scales, and shows (eb) the role of fundamental constants such as (=) the speed of light and Plank's constant h in the fundamental lengths scale without (e) invoking the principles of quantum mechanics. Copyright © 2005 Elsevier Ltd. All rights reserved |

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**Chaos, Solitons & Fractals Volume 14, Issue 9, December 2002, Pages 1341–1351 Fractal strings as an alternative justification for El Naschie's Cantorian spacetime and the fine structure constant Carlos Castro, http://dx.doi.org/10.1016/S0960-0779(02)00110-8**

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| (41) Beginning with the most general fractal strings/sprays construction recently expounded in the book by Lapidus and Frankenhuysen, it is shown how the complexified extension of El Naschie's cantorian-fractal spacetime model belongs to (e&eb) a very special class of families of fractal strings/sprays whose scaling ratios are given by (e) suitable pinary (pinary, p prime) powers of the Golden Mean. |
| (42) They then proceed to show why the logarithmic periodicity laws in (eb) Nature are direct physical consequences of (e) the complex dimensions associated with (e&eb) these fractal strings/sprays. |
| (43) Authors proceed with a discussion on quasi-crystals with (e&eb) p-adic internal symmetries, von Neumann's Continuous Geometry, the role of wild topology in (e&eb) fractal strings/sprays, (e&eb) the Banach-Tarski paradox, and tesselations of (e) the hyperbolic plane, quark confinement and (e&eb) the Mersenne-prime hierarchy of (e) bit-string physics in determining the fundamental physical constants in Nature. Fax: +1-914-965-3397 Copyright © 2002 Published by Elsevier Ltd. About ScienceDirectRemote access Shopping cart Contact and support Terms and conditions Privacy policy Cookies are used by this site. For more information, visit the cookies page. Copyright © 2016 Elsevier B.V. or its licensors or contributors. Science Direct ® is a registered trademark of Elsevier B.V. |

**Chaos, Solitons & Fractals Volume 21, Issue 2, July 2004, Pages 261–271 On Cantorian spacetime over number systems with division by zero Jakub Czajko, http://dx.doi.org/10.1016/j.chaos.2003.12.046**

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| (44) Division by zero is proposed in order to open (eb) new horizons for (e) abstract mathematics and physics. |
| (45) Although the formerly forbidden operation may be of (e) little practical value, the feasibility of unrestricted operations is needed for (e) meaningful deployment of linear vector spaces, whose duals are (=) inverse transposes of their primary vector spaces. |
| (46) Both quantum mechanics and (e&eb) pan-geometry of multispatial hyperspace operate on (e&eb) vectors that belong to (e&eb) mutually dual linear vector spaces. |
| (47) The division by zero also may remove (e) many years old paradox related to (e&eb) multiplication of zeros and makes it possible to introduce (eb) the concept of bigroup as (=) an additive and multiplicative group at the same time. |
| (48) This operational extension to (e&eb) algebraic groups makes thus (eb) an infinite-dimensional, fractal Cantorian spacetime (=) the preferred, invertible, operationally unrestricted, abstract mathematical infrastructure for physics. Tel.: +1-925-681-2058 Copyright © 2003 Elsevier Science Ltd. All rights reserved. |

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The nontrivial equilibrium is stable.

Some of the examples in which stability analysis is of seminal importance are: These are the system to which the

model becomes applicable

**Observation of a Dissipation-Induced Classical to Quantum Transition J. Raftery, D. Sadri, S. Schmidt, H. E. Türeci, and A. A. Houck Phys. Rev. X 4, 031043 – Published 8 September 2014**

(531)Here, authors report the experimental observation of (e) a dynamical quantum phase transition in (eb) a strongly interacting open photonic system.

(532) The system studied, comprising a Jaynes-Cummings dimer realized on (eb) a superconducting circuit platform, exhibits (eb) a dissipation-driven localization transition.

(533)Signatures of the transition in the homodyne signal and photon number reveal (eb) this transition to be from (e) a regime of classical oscillations into (e&eb) a macroscopically self-trapped state manifesting revivals, (=) a fundamentally quantum phenomenon.

(534) This experiment also demonstrates (eb) a small-scale realization of a new class of quantum simulator, whose well-controlled (e&eb) coherent and dissipative dynamics is suited to (e) the study of quantum many-body phenomena out of (e) equilibrium. Received 10 December 2013 DOI:http://dx.doi.org/10.1103/PhysRevX.4.031043 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**Viewpoint: Fermionic Vortices Find their Dual John McGreevy, Department of Physics, University of California at San Diego, La Jolla, CA 92093-0319, USA June 27, 2016• Physics 9, 72Theoretical work reveals a surprising relationship between the physics of fermionic vortices and quantum electrodynamics. Figure caption expand figure APS/Alan Stonebraker Figure 1: The action of discrete symmetries on the fermionic particle (top) and its dual fermionic vortex (bottom).**

(535)Dirac fermions in 2+1 dimensions (two spatial dimensions plus time) have been found in (eb) graphene and on (e&eb) the surfaces of topological insulators (TIs).

(536) A trio of papers now shows that these particles can also be found in (eb) a dual description of quantum electrodynamics (QED) [1–3].

(537)This deceptively simple statement joins (e&eb) together several disparate lines of progress in theoretical physics: interacting topological insulators, (e&eb) anomalies in quantum field theory, (e&eb) spin liquids, and (e&eb) compressible Quantum Hall States.

(538)Humans are not very good at understanding interacting many-body systems.

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| (539)We understand the behavior of free particles and can sometimes use (e) perturbation theory to (e) incorporate weak interactions between (e&eb) them. |
| (540)But what happens when we couple (e&eb) the particles strongly? |
| (541)To reveal the interesting collective phenomena that can result, we've had to rely largely on (e) experiments. |
| (542)Examples include (e) superconductivity (e&eb) superfluidity (e&eb) the integer and (e&eb) fractional quantum Hall effects (e&eb) the fractionalization of charge and spin, and (e&eb) the wetness of water. |
| (543)The fractional quantum Hall effect (FQHE) is a physical phenomenon in which the Hall conductance of 2D electrons shows (eb) precisely quantised plateaus at (eb) fractional values of (e) {\displaystyle e^{2}/h} e^{2}/h. |
| (544) It is a property of (e) a collective state in which electrons bind (e&eb) magnetic flux lines to make (eb) new quasiparticles, and excitations have (e) a fractional elementary charge and possibly also fractional statistics. The 1998 Nobel Prize in Physics was awarded to Robert Laughlin, Horst Störmer, and Daniel Tsui for the discovery and explanation of the fractional Hall effect.[1] |
| (545) However, Laughlin's explanation was (=) a phenomenological guess and only applies to (e&eb) fillings {\displaystyle \nu =1/m} \nu =1/m where (e) {\displaystyle m} m is an odd integer. The microscopic origin of the FQHE remains unknown and is a major research topic in condensed matter physics. (Wikipedia) |
| (546)The quantum Hall effect (or integer quantum Hall effect) is a quantum-mechanical version of (e) the Hall effect, observed in (eb) two-dimensional electron systems subjected to (e&eb) low temperatures and strong magnetic fields, in which the Hall conductance σ undergoes (e&eb) certain quantum Hall transitions to take on (e&eb) the quantized values {\displaystyle \sigma ={\frac {I_{\text{channel}}}{V_{\text{Hall}}}}=\nu \;{\frac {e^{2}}{h}},} \sigma = \frac{I_\text{channel}}{V_\text{Hall}} = \nu \; \frac{e^2}{h}, where {\displaystyle I_{\text{channel}}} I_\text{channel} is the channel current, {\displaystyle V_{\text{Hall}}} V_\text{Hall} is the Hall voltage, e is the elementary charge and h is Planck's constant. The prefactor ν is known as the "filling factor", and can take on either integer (ν = 1, 2, 3, ...) or fractional (ν = 1/3, 2/5, 3/7, 2/3, 3/5, 1/5, 2/9, 3/13, 5/2, 12/5, ...) values. |
| (547) The quantum Hall Effect is referred to as (=) the integer or fractional quantum Hall effect depending on whether ν is an integer or fraction, respectively. |
| (548)The striking feature of the integer quantum Hall Effect is (=) the persistence of the quantization (i.e. the Hall plateau) as (e) the electron density is varied. |
| (549) Since the latter remains constant when (e) the Fermi level is in a clean spectral gap, it must be that the Fermi level is immersed in (eb) essential spectrum (i.e. of finite density of states), though Anderson localized. |
| (550) The quantization and stability of the Hall conductance in such conditions was achieved in 1994 by (e) Bellissard, van Elst and Schulz-Baldes using the methods of (e) Non-Commutative Geometry.[1] |
| (551)The fractional quantum Hall effect is more complicated, as its existence relies fundamentally on (e) electron–electron interactions. |
| (552)Although the microscopic origins of the fractional quantum Hall effect are unknown, there are several phenomenological approaches that provide (eb) accurate approximations. |
| (553) For example the effect can be thought of as (=) an integer quantum Hall effect, not of electrons but of (e) charge-flux composites known as (=) composite fermions. |

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| (554) In 1988, it was proposed that there was quantum Hall effect without Landau levels. This quantum Hall Effect is referred to as (=) the quantum anomalous Hall (QAH) effect. |
| (555)There is also a new concept of the quantum spin Hall effect which is an analogue of the quantum Hall effect, where (e) spin currents flow instead of (e&eb) charge currents.[2] |
| (556)In two dimensions, when classical electrons are subjected to (e&eb) a magnetic field they follow (eb) circular cyclotron orbits. |
| (557)When the system is treated quantum mechanically (eb) these orbits are quantized. |
| (558)The energy levels of these quantized orbitals take on (e) discrete values |
| (559):{\displaystyle \textstyle E_{n}=\hbar \omega _{c}(n+1/2),} \textstyle E_{n}=\hbar \omega _{c}(n+1/2), where ωc = eB/m is the cyclotron frequency |
| (560) These orbitals are known as Landau levels, and at weak magnetic fields, their existence gives rise to (eb) many interesting "quantum oscillations" such as the Shubnikov–de Haas oscillations and the de Haas–van Alphen effect (which is often used to map the Fermi surface of metals). |
| (561) For strong magnetic fields, each Landau level is (=) highly degenerate (i.e. there are many single particle states which have the same energy En). |
| (562)Specifically, for a sample of area A, in magnetic field B, the degeneracy of each Landau level is (=) {\displaystyle \textstyle N=g_{s} BA/\phi _ {0},} \textstyle N=g_{s} BA/\phi _ {0}, where gs represents (eb) a factor of 2 for spin degeneracy, and ϕ0 = 2·10−15 Wb is the magnetic flux quantum. |
| (563) For sufficiently strong B-fields, each Landau level may have (e) so many states that all of the free electrons in the system sit in only a few Landau levels; it is in this regime where (e) one observes the quantum Hall Effect. |
| (564)The integers that appear in the Hall Effect are (=) examples of **topological quantum numbers**. |
| (565) They are known in mathematics as the first **Chern numbers** and are closely related to (e&eb) **Berry's phase.** |
| (566)A striking model of much interest in this context is the Azbel-Harper-Hofstadter model whose quantum phase diagram is (=) the Hofstadter butterfly shown in the figure. (deleted due to spatial constraints) |
| (567)The vertical axis is (=) the strength of the magnetic field and the horizontal axis is the chemical potential, which fixes the electron density. The colors represent the integer Hall conductance. Warm colors represent positive integers and cold colors negative integers. The phase diagram is fractal and has structure on all scales. In the figure there is an obvious self-similarity. |
| (568)Concerning physical mechanisms, impurities and/or particular states (e.g., edge currents) are important for (e) both the 'integer' and 'fractional' effects. |
| (569) In addition, Coulomb interaction is also essential in (eb) the fractional quantum Hall effect. |
| (570)The observed strong similarity between integer and (=) fractional quantum Hall effects is explained by (e) the tendency of electrons to form bound states with (e&eb) an even number of magnetic flux quanta, called composite fermions. |
| (571)Duality is the circumstance in which cranking up interactions results in (eb) another system we understand, usually one that can be described in terms of (e&eb) another set of weakly coupled particles. |
| (572)When such a dual representation is found, (eb) it gives a solution to the strong coupling problem. |
| (573) A further reason that theoretical physicists get excited about dualities is that they undermine (e) the very notion of an “elementary particle.” |
| (574)In terms of the original variables, the dual, weakly coupled particles are (=) generally solitons—large, heavy, classical objects made up of (e) many quanta. |

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| (575)The remarkable thing is that solitons can become (=) "elementary," in the sense that the model that describes the system can be reformulated with the solitons as (=) the basic constituents. |
| (576)An example of a duality, which we will use below, arises in (eb) Maxwell's theory of electromagnetism, where (e) the strength of the interactions is determined by (e) the electric charge e. |
| (577)The duality exchanges electric and (e&eb) magnetic point charges, and the interaction strength in the dual description is 1⁄e, which is small when (e) e is large. |
| (578)This idea has been generalized to (e) many other field theories, thanks, in particular, to tools of (e) supersymmetry [4]. |
| (579)Another classic example of duality involves (e&eb) vortices. |
| (580) A vortex is a localized object in the plane around which a phase field winds (e&eb). (A phase field takes values on the circle.) |
| (581)Dualities relating particles and vortices have (e) a rich history in particle physics and condensed-matter physics. |
| (582) The simplest version relates the particles in the normal state of a superfluid to (e&eb) the vortices in a dual superconductor [5]. |
| (583)This charge-vortex duality has been (=) a workhorse in the study of two-dimensional condensed-matter systems, such as (=) fractional Quantum Hall States, superconducting thin films, and beyond-Landau critical phenomena [6]. |
| (584)Vortices can be (=) fermions (as they frequently are, for instance, in type-II superconductors). |
| (585) A dual description of a system with (e&eb) fermionic vortices thus requires (e) a new ingredient, since (e) it must include (e) such a fermionic particle. |
| (586)The new work [1–3] builds on (e&eb) earlier theoretical work [7] to make (eb) an appealing proposal for (e) the required duality: the fermionic vortex of QED in (eb) 2+1 dimensions is (=) a free Dirac fermion. The argument for this proposal builds on our understanding of TIs. |
| (587)TIs are distinguished only by (e) their edge states. |
| (588) The most famous example is a (3+1)-dimensional time-reversal invariant TI, which can host boundary states whose band structure is defined by (e) a single Dirac cone. |
| (589)A single time-reversal-symmetric Dirac cone is not possible without (e) the bulk TI, so its appearance in experiments provides (eb) a signature of the bulk phase. |
| (590) But the Dirac cone is not (e) the only possible signature. |
| (591) Deformations of the boundary conditions can dramatically change (e&eb) the spectrum of the edge states, yet preserve some essential "TI-ness." This essential property is (=) called an anomaly in high-energy theory. |
| (592)Max Metlitski of the University of California, Santa Barbara, and Ashvin Vishwanath of the University of California, Berkeley [1], and independently, Chong Wang and Todadri Senthil at the Massachusetts Institute of Technology, Cambridge [2], use this ambiguity in the edge theory and the electric-magnetic duality of an auxiliary theory of bulk electromagnetism coupled to (e&eb) the TI to argue for the duality. |
| (593)When I said that (3+1)-dimensional QED had a duality that interchanged the electric and magnetic charges, one could have (e) complained that the electron is (=) a fermion, while (e) a magnetic monopole may not be. |
| (594)However, the authors observe that in a TI, a charge-two monopole is (=) also a fermion. |

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| (595) By analyzing the spectrum of electric and magnetic charges, they conclude that the duality exchanges the electron and (e&eb) the charge-two monopole and, moreover, exchanges a TI with (e&eb) time-reversal symmetry with (e&eb) a "chiral TI," which instead has (e) particle-hole symmetry |
| (596). But this is just what is required of (e) a symmetry that exchanges a fermionic particle with (e&eb) a fermionic vortex: time-reversal symmetry maps the electron to itself but maps the vortex to (e&eb) another vortex whose phase winds (e&eb) the other way (see Fig. 1). |
| (597)The exchange of particles and holes, in turn, preserves (eb) the vortex. |
| (598)The bulk duality implies (eb) a duality between edge theories. |
| (599)The freedom to deform the edge of the TI implies (eb) a weak form of equivalence: although the edge theories are (=) topologically the same (in terms of symmetry assignments and (e&eb) anomalies), they need not have (e) identical low-energy physics. |
| (600)One side of the duality is (=) a free, massless Dirac particle, but the long-distance physics of (2+1)-dimensional QED is (=) a long-standing problem, which is important in the study of spin liquids. |
| (601) Many gauge theories in two dimensions develop (eb) a gap in their energy spectrum because of (e) the proliferation of monopoles. |
| (602) In the realization proposed by Metlitski and Vishwanath [1] and by Wang and Senthil [2], this cannot happen because (e) the monopoles carry electric charge, and indeed the stronger notion of duality as (=) an equivalence of low-energy physics may hold. |
| (603)Evidence for the stronger relation comes from (e) David Mross and colleagues at the California Institute of Technology, Pasadena [3], who have derived a lattice version of the duality using (e) a clever but explicit change of variables. |
| (604)Their construction is particularly interesting because (e) it uses (e) a lattice of one-dimensional models which themselves can only occur at the edge of a TI. |
| (605)Authors explicitly derive the duality between a free electronic Dirac cone and (e&eb) quantum electrodynamics in (eb) (2+1) dimensions (QED3) with (e&eb) N=1 fermion flavors. |
| (606)The duality proceeds via (e&eb) an exact, nonlocal mapping from electrons to (e&eb) dual fermions with (e&eb) long-range interactions encoded by (e) an emergent gauge field. |
| (607)This mapping allows (eb) us to construct (eb) parent Hamiltonians for exotic topological-insulator surface phases, derive (eb) the particle-hole-symmetric field theory of (e) a half-filled Landau level, and nontrivially constrain (e) QED3 scaling dimensions. |
| (608)They similarly establish duality between bosonic topological insulator surfaces and (e&eb) N=2 QED3. |

**Black hole thermodynamics with generalized uncertainty principle Li Xiang, X. Q. Wen (Submitted on 6 Jan 2009 (v1), last revised 18 Mar 2009 (this version, v2))**

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| (609) In the standard viewpoint, the temperature of a stationary black hole is proportional to (=) its surface gravity, $T_H=\hbar\kappa/2\pi$. |
| (610)This is a semiclassical result and the quantum gravity effects are not (e) taken into consideration. |
| (611)This Letter explores a unified expression for (e) the black hole temperature in the sense of (e) a generalized uncertainty principle (GUP). |
| (612) Discussion involves a heuristic analysis of (e) a particle which is absorbed by (e) the black hole. |
| (613)Besides a class of static and spherically symmetric black holes, an axially symmetric Kerr-Newman black hole is considered. Different from the existing literature, authors suggest (eb) that the black hole's irreducible mass represent (eb) the characteristic size in (eb) the absorption process. |

(614)The information capacity of a remnant is also discussed by (e) Bousso's D-bound in de Sitter spacetime. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP 0910:046,2009 DOI: 10.1088/1126-6708/2009/10/046 Cite as: arXiv:0901.0603 [gr-qc] (or arXiv:0901.0603v2 [gr-qc] for this version)

IF we define

$$[a_7]_7 = [f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}]\,P_7^*$$

$$[a_7]_7 = \left[f_{7*}^{7}\right] P_7^* - \frac{h_*^7}{h_*} P^*(0)$$

$$e = [g_*^7]_7\, Q_7^*$$

$$\eta_7(t) = \left[h_*\right]_7 \beta_7 \left(\theta_7(t)\right) - \mu_7(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\left[a_7\right]M_7 - [a_7]_7 W_7 - [n(t)]M_7$$

$$w_7(t) = eM_7$$

Write

$$A = \begin{pmatrix} -[a_7]_7 - [a_7]_7 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)= \begin{pmatrix} [\eta(t)]_7 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_7 \\ W_7 \end{pmatrix}$$

The solution of (32)is

$$X_7(t) = e^{A,t}[X]_7(0) + \int_0 e^{At}$$

## PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION

## ATTRITION MODELS

## MODULE EIGHT

Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal.

### 52. MODEL FORMULATION

**Model is applied to systems given in the following:**

**Researchers blur the line between classical and quantum physics by connecting chaos and entanglement July 12, 2016 by Sonia Fernandez**

(3120) Experimental link between quantum entanglement (left) and (e&eb) classical chaos (right) found using (e) a small quantum computer. Credit: University of California - Santa Barbara Using a small

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| quantum system consisting of three superconducting qubits, researchers at UC Santa Barbara and Google have uncovered a link between aspects of classical and (e&eb) quantum physics thought to be unrelated: classical chaos and (e&eb) quantum entanglement. |
| (3121) Their findings suggest (eb) that it would be possible to use (e) controllable quantum systems to investigate certain fundamental aspects of nature. |
| (3122) "It's kind of surprising because (e) chaos is this totally classical concept—there's (=) no idea of chaos in a quantum system," Charles Neill, a researcher in the UCSB Department of Physics and lead author of a paper that appears in Nature Physics. |
| (3123) "Similarly, there's no concept of entanglement within (eb) classical systems. And yet it turns out that chaos and (e&eb) entanglement are really very strongly and clearly related." |
| (3124) Initiated in the 15th century, classical physics generally examines and describes (eb) systems larger than atoms and molecules. |
| (3125) It consists of hundreds of years' worth of study including Newton's laws of motion, electrodynamics, relativity, thermodynamics as well as chaos theory—the field that studies the behavior of (e) highly sensitive and unpredictable systems. |
| (3126) One classic example of chaos theory is the weather, in which a relatively small change in one part of (e) the system is enough to foil predictions—and vacation plans—anywhere on the globe. |
| (3127) At smaller size and length scales in nature, however, such as those involving (e&eb) atoms and photons and their behaviors, classical physics falls short. |
| (3128) In the early 20th century quantum physics emerged, with its seemingly counterintuitive and sometimes controversial science, including (e) the notions of superposition (the theory that a particle can be located in (eb) several places at once) and entanglement (particles that are deeply linked behave as (=) such despite physical distance from one another). |
| (3129) And so began the continuing search for connections between the two fields. All systems are fundamentally quantum systems, according Neill, but the means of describing in a quantum sense the chaotic behavior of, (e) say, air molecules in an evacuated room, remains limited. |
| (3130) Imagine taking a balloon full of air molecules, somehow tagging them so you could see them and then releasing them (eb) into a room with no air molecules, noted co-author and UCSB/Google researcher Pedram Roushan. |
| (3131) One possible outcome is that the air molecules remain clumped together (e&eb) in a little cloud following the same trajectory around the room |
| (3132) And yet, he continued, as we can probably intuit, the molecules will more likely take off in (eb) a variety of velocities and directions, bouncing off walls and interacting with (e&eb) each other, resting after the room is (=) sufficiently saturated with them. |
| (3133) "The underlying physics is (=) chaos, essentially," he said. |
| (3134) The molecules coming to rest—at least on (eb) the macroscopic level—is the result of (e) thermalization, or of reaching equilibrium after they have achieved (eb) uniform saturation within the system. |
| (3135) But in the infinitesimal world of quantum physics, there is (=) still little to describe that behavior. |
| (3136) The mathematics of quantum mechanics, Roushan said, do not allow for (e) the chaos described by Newtonian laws of motion. |
| (3137) To investigate, the researchers devised an experiment using (e) three quantum bits, the basic computational units of (e) the quantum computer. |

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| (3138) Unlike classical computer bits, which utilize (e) a binary system of two possible states (e.g., zero/one), a qubit can also use (e) a superposition of both states (zero and one) as a single state. |
| (3139) Additionally, multiple qubits can entangle, or link so closely that their measurements will (=) automatically correlate. |
| (3140) By manipulating these qubits with electronic pulses, Neill caused them to interact, rotate and evolve (e&eb, e&eb, eb) in the quantum analog of (e) a highly sensitive classical system. |
| (3141) The result is a map of entanglement entropy of (e) a qubit that, over time, comes to strongly resemble that of (e) classical dynamics—the regions of entanglement in (eb) the quantum map resemble the regions of (e) chaos on the classical map. |
| (3142) The islands of low entanglement in (eb) the quantum map are located in the places of (e) low chaos on the classical map. |
| (3143) "There's a very clear connection between entanglement and (e&eb) chaos in these two pictures," said Neill. |
| (3144) "And, it turns out that thermalization is (=) the thing that connects chaos and (e&eb) entanglement. |
| (3145) It turns out that they are actually the driving forces behind (eb) thermalization. |
| (3146) "What we realize is that in almost any quantum system, including on quantum computers, if you just let it evolve and you start to study what happens as a function of time, (eb) it's going to thermalize," added Neill, referring to the quantum-level equilibration. "And this really ties together the intuition between classical thermalization and (e&eb) chaos and how it occurs (eb) in quantum systems that entangle." |
| (3147) The study's findings have fundamental implications for (e) quantum computing. |
| (3148) At the level of three qubits, the computation is relatively simple, said Roushan, but as researchers push to build increasingly sophisticated and powerful quantum computers that incorporate (e) more qubits to study highly complex problems that are beyond the ability of classical computing—such as those in the realms of machine learning, artificial intelligence, fluid dynamics or chemistry—a quantum processor optimized for such calculations will be a very powerful tool. |
| (3149) "It means we can study things that are completely impossible to study right now, once we get to bigger systems," said Neill.More information: Nature Physics, DOI: 10.1038/nphys3830 , http://www.nature.com/nphys/journal/vaop/ncurrent/full/nphys3830.html Journal reference: Nature Physics search and more info website Provided by: University of California - Santa Barbara |

**Title: coherent control of macroscopic quantum states in a single-Cooper-pair box Authors: Nakamura, Y.; Pashkin, Yu. A.; Tsai, J. S. Publication: Nature, Volume 398, Issue 6730, pp. 786-788 (1999). (Nature Homepage) Publication Date: 04/1999 Origin: NATURE DOI: 10.1038/19718 Bibliographic Code: 1999Natur.398...786N**

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| (3150) A nanometre-scale superconducting electrode connected to (e&eb) a reservoir via (e&eb) a Josephson junction constitutes (e) an artificial two-level electronic system: a single-Cooper-pair box. |
| (3151) The two levels consist of (e) charge states (differing by 2e, where (e) e is the electronic charge) that are coupled by (e) tunnelling of Cooper pairs through (e&eb) the junction. |
| (3152) Although the two-level system is macroscopic, containing (e) a large number of electrons, the two charge states can be (=) coherently superposed. |

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| (3153) | The Cooper-pair box has therefore been suggested as (=) a candidate for a quantum bit or `qubit'- the basic component of a quantum computer. |
| (3154) | Here authors report the observation of quantum oscillations in (eb) a single-Cooper-pair box. |
| (3155) | By applying a short voltage pulse via (e&eb) a gate electrode, we can control (e&eb) the coherent quantum state evolution: the pulse modifies (e&eb) the energies of (e) the two charge states non-adiabatically, bringing them into (e&eb) resonance. |
| (3156) | The resulting state-a superposition of (e) the two charge states-is detected by (e) a tunnelling current through (e&eb) a probe junction. |
| (3157) | Results demonstrate electrical coherent control of (e) a qubit in a solid-state electronic device. |

**Superconducting quantum computing From Wikipedia, the free encyclopedia**

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| (3158) | Superconducting quantum computing is a promising implementation of (e) quantum information technology that involves (e&eb) nanofabricated superconducting electrodes coupled through (e&eb) Josephson junctions. |
| (3159) | As in a superconducting electrode, the phase and (e&eb) the charge are conjugate variables. |
| (3160) | There exist three families of superconducting qubits, depending on (=) whether the charge, (e&eb) the phase, or neither of (e) the two is good quantum number. |
| (3161) | These are respectively termed charge qubits (e&eb) flux qubits, and (e&eb) hybrid qubits. |

**Theory**

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| (3162) | Unlike many other physical implementations of a qubit which involve (e&eb) exclusively two-level systems (such as nuclear spin and (e&eb) photon polarization), the integrated quantum circuit involved in (e&eb) a superconducting qubit is (=) a multi-level system, of which only the first two levels are used as (=) the computational basis. |
| (3163) | A basic requirement for such an implementation is that the energy levels are (=) not uniformly spaced, so that (e) photons of a particular frequency which cause (eb) transition between the 0 and 1 levels do not(e) cause transitions from the first level to (e&eb) the higher levels as well. |
| (3164) | For electronic signals to be carried from one part of the circuit to (e&eb) another without energy loss and (e&eb) hence decoherence, the system also needs to be (=) non-dissipative, i.e., the metallic parts involved should have (e) zero resistance. |
| (3165) | These circuits currently need to be (=) operated at very low temperatures so that, 1) superconductivity is (=) realized, and 2) thermal fluctuations do not (e) cause transitions between (e&eb) energy levels. |

**Quantum LC Circuit**

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| (3166) | The simplest non-dissipative quantum circuit consists simply of (e) an inductor L {\displaystyle L} L and (e&eb) capacitor C {\displaystyle C} C, with (e&eb) the metallic wires connecting them being superconducting. |
| (3167) | Here the flux Φ {\displaystyle \Phi } \Phi through (e&eb) the inductor and (e&eb) the charge Q {\displaystyle Q} Q in the capacitor are (=) canonically conjugate variables which obey (eb) the commutation relation [ Φ , Q ] = i ℏ {\displaystyle [\Phi ,Q]=i\hbar } [\Phi ,Q]=i\hbar |

**The Hamiltonian associated with this system is**

(3168) H = Φ 2 2 L + Q 2 2 C {\displaystyle H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}} H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}

(3169) This gives rise to the energy levels E = ℏ ω 0 ( n + 1 2 ) {\displaystyle E=\hbar \omega _{0}(n+{\frac {1}{2}})} E=\hbar \omega _{0}(n+{\frac {1}{2}}) Here ω 0 = 1 L C {\displaystyle \omega _{0}={\frac {1}{\sqrt {LC}}}} \omega _{0}={\frac {1}{{\sqrt {LC}}}}

**Quantum gravity Carlo Rovelli (2008), Scholarpedia, 3(5):7117. doi:10.4249/scholarpedia.7117 revision #135539 [link to/cite this article] Post-publication activity Curator: Carlo Rovelli Dr. Carlo Rovelli, Center of Theoretical Physics, and Marseille, France**

(3170) Quantum Gravity is the name given to any theory that describes gravity in (eb) the regimes where quantum effects cannot be disregarded.

(3171) At present, there is no such a theory which is universally accepted and confirmed by experience. Therefore the term "Quantum Gravity" indicates (eb) more an open problem than a specific theory.

(3172) Several research lines, at different levels of development, offer tentative solutions to the problem. These tentative quantum-gravity theories are variously viewed as competing research directions, or as contributions to the common goal of finding the physically correct theory. The quest for the good quantum theory of gravity bears on a number of fundamental issues and it is sometimes presented as the most important open problem in fundamental physics: the "Holy Grail" of contemporary theoretical physics.

**The problem of quantum gravity**

(3173) The current theory of (e) gravity is general relativity, a theory developed by Albert Einstein at the beginning of the XX century, which has changed in depth our understanding of (e) what are space and time.

(3174) Similarly, quantum mechanics has changed in depth our understanding of matter, energy and causality. These two theories, today very well empirically supported, have opened (eb) a major conceptual revolution in physics.

(3175) However, they appear to be incompatible, at least at a first reading, because (e) each of the two is formulated on the basis of (e) principles that are explicitly contradicted by (e) the other theory.

(3176) Therefore in spite of the immense development of scientific knowledge it has delivered, the physics of the XX century has left us with (e&eb) a major confusion on what are the basic conceptual ingredients for (e) understanding the physical world.

(3177) In other words, the scientific revolution opened by general relativity and quantum mechanics at the beginning of (e) the XX century is not concluded yet, and a new synthesis is required. Quantum gravity, merging general relativity and quantum mechanics should be this new synthesis.

**Quantum space and quantum time**

(3178) Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.

(3179) In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.

(3180) They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.

(3181) When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.

(3182) Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.

**The Planck scale**

(3183) Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".

(3184) Here $\hbar$ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.

(3185) The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.

(3186) Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)

(3187) Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.

## Physical applications

Some phenomena where quantum gravity plays a major role are the following.

**The microscopic structure of spacetime**

(3188) If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.

(3189) Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length

(3190) Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.

(3191) A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.

**Early cosmology**

(3192) According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.

(3193) However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.

(3194) A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.

(3195) Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".

(3196) One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.

**Black holes**

(3197) Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.

(3198) First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.

(3199) Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.

(3200) What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.

(3201) Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.

**Astrophysical effects**

(3202) Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.

(3203) An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.

(3204) The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.

**Main approaches**

(3205) The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.

(3206) Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.

(3207) The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.

**String theory**

(3208) String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.

(3209) Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.

(3210) String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.

(3211) Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.

(3212) However, when summed the entire series appears to be (=) divergent.

(3213) In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.

(3214) Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.

(3215) Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.

(3216) In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.

**Loop quantum gravity and spin foams**

(3217) Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.

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| (3218) The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity. |
| (3219) Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime. |
| (3220) The granular properties of space can be (=) explicitly computed. |
| (3221) In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory. |
| (3222) The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks. |
| (3223) Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure. |
| (3224) The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit. |
| (3225) Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism. |
| (3226) This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism. |

**Noncommutative geometry**

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| (3227) Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime. |
| (3228) Quantum mechanics teaches us that dynamical quantities are (=) noncommutative --meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed. |
| (3229) It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry. |
| (3230) A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces. |
| (3231) Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces. |
| (3232) This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity. |

**Others**

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| (3233) Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays |

(3234) . Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection.

(3235) Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible.

(3236) The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.

Let $p_8(t, a_8)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(g_8)Q_8$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_8(0) = 0, \quad g_8(P_8) > 0, \quad P_8 \; 0, \qquad \lim_{P_8 \to \infty} g_8(P_8) = (K_8)_8 \; \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_8 + (\rho_a)_8 + D_8(a_8, P_8, Q_8)\rho_8 + 0, \qquad 0 < a_8 \; \infty,$$

$$\rho_8(t,0) = \int_0^\infty n_8(a_8, P_8, Q_8)\rho_8(t_8, a_8)da_8$$

Where $D_8$ is the mortality function and $n_8$ is the fertility function.

Put $_8$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_8}{dt}\right]_8 = x_8\left[\frac{dQ_8}{dt}\right]_8 - x_8\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_8}{dt}\right]_8 = x_3\left[\frac{dQ_3}{dt}\right]_8 - x_4\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_8}{dt}\right]_8 - x_6\left[\frac{dQ_3}{dt}\right]_8$$

Where $x_8, x_8, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_8$ and $g_8(P_8)Q_8$

Now when you add the term $(b_8Q_8)$ with $g_8(P_8)(Q_8)$ we are infect see is that $(b_8Q_8)$ is credited incrementally part of $g_8(P_8)(Q_8)$ or instant nearly, when subtracted from $\frac{dQ_8}{dt}$ , it means that $g_8(P_8)(Q_8)$ is slowly eaten by $\frac{dQ_8}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_8 = \mu_8(a_8) + (f_8)_8(P_8)_8 + (f_8)_8(Q)_8 \quad (2)$$

$$n_8 = \beta_8(a_8)h_8(Q_8), \quad (3)$$

Where $\mu_8$the age- dependent debit rate whit is $\mu_8(a_8)$ 0 $\forall a > 0$; $(f_8)_8$ is the density- dependent death rate with $f_8(0) = 0$, $(f_8^8)_8(P_8)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_8)_8(P_8) = (k_8)_8$ ; $f_8$ is the functional response of creditors to prey with $[f_8]_8(0) = 0, (f_8^8)_8(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_8)_8(Q_8) = [k_8]$ ; $\beta$ is the age

dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^8(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_8 + (\rho_a)_8 = -\mu_8(a_8) + (f_8)_8 P_8) + \big(f_8)_8(Q_8)\big)(\rho_8)_8 \quad (4)$$

$$\rho_8(t_8, 0) = h_8(Q_8) \int_0^\infty \beta_8 (a_8)\rho_8(t, a_8)da, \quad (5)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(P_8)Q_8 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_8$ and $Q_8$ of (4-6) satisfy

$$0 \ \rho_8(t, a_8) \ \rho_8^*(t, a_8),$$

$$0 \ Q_8(t) \ Q_8^*(t),$$

Where $\rho_8^*(t, a_8)$ and $Q_8^*(t)$, are solutions of

$$(\rho_t^*)_8 + (\rho_a^*)_8 = -(\mu_8(a_8) + (f_8)_8(P_8^*) + (f_8)_8 Q_8^*))\ \rho_8^*, \quad (7)$$

$$\rho_8^*(t, 0) = \int_0^\infty \beta_8(a_8)\ \rho_8^*(t, a_8)da_8 \quad (8)$$

$$\frac{dQ_8^*}{dt} = -b_8\ Q_8^* + g_8(P_8^*)Q_8^*, \quad (9)$$

With $P_8^*(t) = \int_0^\infty \rho_8^*\ (t, a_8)da_8$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_8^*(t)}{dt} = P_8^*(t)\big(r_8 - (f_8)(P_8^*) - (f_8)_8(Q^*)\big), \tag{10}$$

$$\frac{dQ_8^*(t)}{dt} = Q_8^*(t)(-b_8 + g_8(P_8^*)), \tag{11}$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_8(a)e^{-\int_0^a \mu_8(\tau)dt}e^{-ra}\ da_8 = 1. \tag{12}$$

If $r < 0$, the trivial solution $P_8^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(t)dt}\ da_8 > 1, \tag{$H1$}$$

Which guarantees $r_8 > 0$.

*Remark.* If $\beta_8(a) = [\beta_0]_8 a_8 e^{-\alpha_8 a_8}, \mu_8(a_8) \equiv \mu_8$ then the solution of (12) is

$$r_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_8]_8(P_8) = 0, [f_8](Q_8) = \lambda_8 Q_8\ h_8(Q_8) = 1, \qquad and\ \ g_8(P_8) = c_8 P_8$$

$$[\rho_8]_8 + [\rho_a]_8 = -(\mu_8 + \lambda_8 Q_8)\rho_8$$

$$\rho_8(t,0) = \int_0^{\infty} [\beta_0]_8 a_8 e^{-\alpha_8 a_8} \rho_8(t,a_8)\, da_8$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + c_8 P_8 Q_8$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_8}{dt} = \omega_8 P_8 - \lambda_8 P_8 Q_8$$

$$\frac{dQ_8}{dt} = b_8 Q_8 - c_8 P_8 Q_8$$

Where $\omega_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$

## 53. THE EQUILIBRIUM $(\rho_0)_8(a_8), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_8(a_8), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_8 + (\rho_a)_8 = -\big(\mu_8(a_8) + (f_8)_8(P_8)\big)\rho,$$

$$\rho_8(t,0) = \int_0^{\infty} \beta_8(a_8)\rho_8(t,a_8)da_8 \qquad (13)$$

The solution of (13) has the form of

$$\rho_8(t,a_8) = Ce_8^{r(t-a_8)} - \int_0^a (\tau)d\tau - \int_0^8 (f_8)_8\big(P_8(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_8(t) = C_8 e^{r_8 t - \int_0^t (f_8)_8(P_8(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_8^{-8}(r)$ if $k_8 > r$. This is equivalent to

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{(k_8)_8 a_8} da_8 < 1.$$

In fact, if we define

$$R_8(s_8) = \int_0^\infty \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau)d\tau} e^{-s_8 a_8}\, da_8$$

$R(s)$ is a decreasing function. $R_8(r_8) = 1$ And $R_8(k_8)_8 < 1$ imply $r_8 < (k_8)_8$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_8(a_8), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{k_8 a_8} da_8 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_8 \equiv \int_0^\infty (\rho_0)_{88}(a_8)da_8$ *satisfies*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{-(f_8)_8(P_0)a_8} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_8(a_8), 0)$

Let $u_8(t,a_8) = \rho(t,a_8) - \rho_0(a_8)$ Then

$$N_8(t) \equiv \int_0^\infty u_8(t,a_8)da_8 = \int_0^\infty \rho_8(t,a_8)da_8 - \int_0^\infty (\rho_0)_8(a_8)da = P_8(t) - P_0)_8$$

And

$$(u_8)_8 + (u_a)_8 = (\rho_t)_8 + (\rho_a)_8 - (\rho_0^8)_8(a_8)$$

$$= -\mu_8(a_8) + (f_8)_8(P_8) + (f_8)_8(Q_8))\,(u_8 + (\rho_0)_8) + \mu_8(a_8)(\rho_0)_8 + (f_8)_8(P_0)_8(\rho_0)_8$$

$$= -\mu_8(a_8)u_8 - (f_8)_8(P_8) + (f_8)_8(Q_8))\,u_8 - (\rho_0)_8\,(f_8)_8(P_8) - (f_8)_8(P_0)_8$$

$$+ (f_8)_8(Q_8)), \qquad (15)$$

Where $(\rho_0^8)_8(a_8)$ denotes $d(\rho_0)_8(a_8)/da_8$

Expansion $(f_8)_8(P_8)$ and $(f_8)_8(Q_8)$ at $(P_0)_8$ and 0 yield respectively

$$(f_8)_8(P_8) = (f_8)_8((P_0)_8 + N_8)$$

$$= (f_8)_8(P_0)_8 + (f_8^8)_8(P_0)_8N_8 + o_8(N_8), \qquad (16)$$

$$(f_8)_8(Q_8)$$

$$= (f_8^8)_8(0)Q + o_8(Q_8) \qquad (17)$$

Write $f_8(P_0) = f_{80}, f_8'(P_0) = f_{80}'$, and $f_8'(0) = f_{80}'$ and substitute (16) and (17) into (15) to obtain

$$(u_8)_8 + (u_a)_8 \approx -(\mu_8(a_8) + (f_{80})_8)u - (f_{80}^8)_8(\rho_0)_8(N_8) - (f_{80}^8)_8(\rho_0)_8(Q_8)$$

$$u_8(t,0) = \rho_8(t,0)(\rho_0)_8(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a) + (\rho_0)_8(a_8)da_8 - (\rho_0)_8\,(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a_8)da_8 + (h_8(Q_8) - 1)(\rho_0)_8\,(0) \qquad (18)$$

Expanding

$$h_8(Q_8) = 1 + (h^8)_8(0)Q_8 + 0(Q_8) \equiv 1 + h_0^{'}Q_8 + 0Q_8$$

And substituting it into (18), we arrive at

$$u_8(t,0) \approx \int_0^{\infty} \beta_8\,(a_8)u_8(t,a_8)da_8 + (h_{a_8}^{'})(\rho_0)_8 Q_8$$

The linearization is the following :

$$[u_8] + [u_8] = -(\mu_8(a_8) + [t_{80}]_8 u_8 - [\alpha_8]_8(a_8)N_8 - [\alpha_8]_8(a_8)Q_8$$

$$u_8(t,0) = \int_0^{\infty} \beta_8(a_8)u_8(t,a_8)da_8 + [\xi_8]Q_8 \qquad (19)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8$$

where

$$[\alpha_8]_8(a_8) = [t_8^8 a]_8[p,o]_8(a_8)d2(a_8) = [f_{80}^8]_8(\rho\, o\,)_8(a_8)\xi = [\mathrm{h\ o}^8]_8[\rho \mathrm{o}]_8(\mathrm{o}) \text{ and } \mathrm{b}_8 = \mathrm{b}_8 - g_8(P_8)$$

If $g_8(P_0) > b$, clearly, $(\rho_0)_8(a_8), 0)$ is unstable and the predator population is persistent. Assume $g_8(P_0)_8 < b$.

Then, $Q_8(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_8 + u_a = -(\mu_8(a_8) + (f_{80})_8)u_8 - (\alpha_8)_8(a_8)N_8$$

$$= -(\mu_8(a_8) + f_{80})u_8 - \alpha_8(a_8)\int_0^{\infty} u_8(t,a_8)da_8$$

$$u_8(t,0) = \int_0^\infty \beta_8\,(a_8)u_8(t,a_8)da_8$$

Using the method developed by Cushing [2, 3], we suppose

$$u_8(t,a_8) = \tilde{u}_8(a_8)^{e^{\gamma(t-a_8)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_8}\,da_8$. Then,

A solution of the first differential equation for $\tilde{u}(a_8)$ gives,

$$\tilde{u}_8(a_8) = e^{-\int_0^a \mu_8(\tau_8)d\tau - (f_{80})_8 a}\left(\tilde{u}_8(0) - w\int_0^a e^{\int_0^\tau \mu_8(s_8)ds + (f_{80})}\tau_8\,(\alpha_8)_8(\tau_8)e^{\gamma_8\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\right)\tilde{u}_8(0)$$

$$+w\int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\int_0^a (\alpha_8)_8(\tau_8)e^{-\gamma_8\tau_8}e^{(f_{80})_8\tau_8 + \int_0^\tau \mu_8(s_8)ds}d\tau_8\,da_8$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}da\right)(\tilde{u}_8(0))_8$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_8 a_8-(f_{80})_8 a_8-\int_0^a \mu_8(\tau_8)d\tau}\int_0^a e^{(f_{80})_8\tau_8+\int_0^\tau \mu_8(s_8)ds+y\tau}\alpha_8(\tau)d\tau\, da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_8 a_8-(f_{80})_8 a_8-\int_0^a \mu_8(t_8)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_8\,(a_8)Y_8(\gamma_8,a_8)da_8\right)u_8(0)$$

$$+w\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)\int_0^a(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8=0,$$

(24)

$$\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\tilde{u}_8(0)$$

$$+\left(-1-\int_0^{\infty}Y_8(\gamma_8,a_8)\int_0^a(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8\right)w=0.$$

Equation (20) has a solution of the form of $u_8(t,a_8)=u_8(a_8)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_8(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_8(\gamma_8,a_8)=\int_0^a(\rho_0)_8(\tau_8)Y^{-8}\,(\gamma_8,a_8)d\tau_8$. We obtain the characteristic equation

$$f_{80_8}^{l}\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\left(\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$+\left(1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)da\right)\left(1+(f_{80}^{l})_8\int_0^{\infty}Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_8(a_8) = (\rho_0)_8(0)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - (f_{80})_8 a_8}$$

$$I_8(\gamma_8, a_8) = \int_0^a (\rho_0)_8(\tau_8)e^{\gamma_8\tau_8 + (f_{80})_8\tau + \int_0^\tau \mu_8(s_8)ds_8}d\tau$$

$$= \frac{(\rho_0)_8(0)}{\gamma_8}(e^{\gamma_8 a_8} - 1)$$

Let

$$E_8(a_8) = e^{-(f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau_8}$$

Then $[Y_8]_8[\gamma_8]_8, a_8) = E_8(a_8)e^{-\gamma_8 a_8}$. by a simple calculation,

$$\int_0^\infty \beta_8(a_8)Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma}\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}\ da\right);$$

$$\int_0^\infty Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma_8}\left(\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}\ da_8\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_8(0)}{\gamma_8}[f_{80}^8]\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}da_8\ (\int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8)$$

$$+\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right).$$

$$\left(1+\frac{(\rho_0)_8(0)}{\gamma_8}(f_{80}^{'})_8\left(\int_0^{\infty}E_8(a_8)da_8-\int_0^{\infty}E_8(a_8)e^{-\gamma_8 a_8}da_8\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right)(1+\frac{(f_{80}^{'})_8}{\gamma_8}(\rho_0)_8(0)\int_0^{\infty}E_8(a_8)da_8)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)da_8=1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\neq 0,$$

Since $\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)\int_0^{a}\alpha_8(\tau_8)Y^{-8}(\gamma_8,\tau_8)d\tau_8 da_8>0$. Therefore, the characteristic solution is

$$\gamma_8=-(f_{80}^{'})_8(\rho_0)_8\int_0^{\infty}E_8(a_8)da_8.$$

Since $f_8^{'}(P_0)_8>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_8),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_8),0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_8 + (\mu_8(a_8) + [f_8]_8(P_8^*) + [f_8]_8(Q_8^*))P_8^* = 0, \quad (26)$$

$$p_8^*(0) = h_8(Q_8^*)\int_0^\infty \beta_8(a_8)p_8^*(a_J)p_8^*(a_8)da_8 \quad (27)$$

$$g_8(P_8^*) - b_8 = 0, \quad (28)$$

Where $P_8^* = \int_0^\infty p_8^*(a_8)\, da_8$

Equation (28) yields $P^* = g^{-8}(b)$,

From (26),

$$p_8^*(a) = p_8^*(0)e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}$$

Substituting into (27), we need to solve

$$\left(1 - h_8(Q_8^*)\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}da_8\right)p_8^*(0) = 0$$

For $Q_8^*$

Fix $p_8^*(0)$ and $P_8^*$ Let

$$H_8(z_8) = 1 - h_8(z_8)\int_0^\infty \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}da_{,8}$$

And assume that

$$[k_3]_8 \int_0^{\infty} \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(g_8^{-8}(b_8))a_8 - [k_8]_8 a_8}\, da_8 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty} [H_8(z_8)] > 0.$$

Clearly,

$$\frac{dH_8}{dz_8} = h_8(z)(\,[f_8^{'}]_8(z_8)\,) \int_0^{\infty} a_8\beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - f_8(P_8^*)a_8 - f_8(z_8)a_8} da_8$$

$-h^{'}(z_8)\int_0^{\infty} \beta_8(a_8)\, e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}\ da_8$

By the assumptions on $\left[f_8^8\right]_8 and\ [h^8]^8 dH_8/dz > 0$. Hence, if

$$H_8(0) = 1 - \int_0^{\infty} \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8} da_8 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_8(0) < 0$ follows if $g_8^-(b_8) < [f_8^{-8}]_8(r_8)$, where $r_8$ is the solution of (12).

Indeed, since

$$R_8(s_8) = \int_0^{\infty} \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 e^{-y_8^a}} da_8$$

Is a decreasing function of $s_8$ when $P_8^* = [g^{-8}]_8(b_8) < [f_8^{-8}]_8(r_8)$, that is $[f_8]_8(P_8^*) < r$, $R\big(f_8(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_8(g^{-8}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8\left(g_8^{-8}(b)\right)a_8} > 1. \qquad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_8^*(a_8), Q_8^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_8^*(a_8), Q_8^*)$ is more complicated. We again consider perturbations $m_8(t_8 a_8) = p_8(t_8, a_8) - p_8^*(a_8)$ and $W_8(t) = Q_8(t) - Q^*$.Then

$M_8(t) \equiv \int_0^\infty m_8(t, a_8)da_8 = P_8(t) - P_8^*,$

$$[m_8]_8 + [m_a]_8 = [p_8]_8 + [p_a]_8 - \left[p^*\right]_8(a_8) \approx -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8$$

$$-[f_{8*}^8]_8 p_8^*(a_8)M - [f_{8*}^8]_8 p_8^*(\mathrm{a})W_8,$$

$$m_8(t,0)p_8(t,0) - p_8^*(0) \approx h_* \int_0^\infty \beta_8(a_8)m_8(t,a_8)\,da + \frac{[h_*^8]_8}{[h_*]_8}\, p_8^*(0)W_8$$

$$\frac{dW_8}{dt} = \frac{dQ_8}{dt} \approx [g_*^8]_8 Q_8^* M_{,8}$$

Where $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $\left[f_{8*}^{'}\right]_8$ denoting $\left[f^{'}{}_8(P^*)\right]_8\left[f_{8*}^{'}\right]_8$ denoting $[f_8]_8$ $(Q_8^*,$ $h_{*}^{'}$ denoting $h_*^{'}$ denoting $h_8^{'}$ $(Q_8^*)$, and $\left[g_*^{'}\right]_8$

Denoting $\left[\dot{g}\right]_8(P_8^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_8]_8 + [m_a]_8 = -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8 - \left[\dot{f}_{8*}\right]_8[p_8^*](a_8)M_8 - \left[\dot{f_{8*}}\right]_8 p_8^*(a_8)W_8$$

$$m_8(t,0) = [h_*]_8 \int_0^\infty \beta_8(a_8)m_8(t,a_8)da_8 + \frac{[h_*^8]_8}{[h^*]_8}[p^*]_8(0)\,W_{,8} \qquad (30)$$

$$\frac{dW_8}{dt} = [g_*^8]_8[Q_8^*]M_8$$

Integrating $[m_8]_8(t,a_8)$ with respect to a from 0 to ∞, we have

$$\frac{dM_8(t)}{dt} = \int_0^\infty [mt]_8(t,a_8)da_8$$

$$= -\int_0^\infty [m_a]_8\,(t,a_8)da_8 - \int_0^\infty \mu(a_8)m(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + \left[\dot{f}_{8*}\right]_8 \int_0^\infty p_8^*(a_8)da_8\right)M_8 - [f_{8*}^8]_8 \int_0^\infty p_8^*(a_8)da_8 W_8$$

$$= m_8(t,0) - \int_0^\infty \mu_8(a_8)\,m_8(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + \left[\dot{f_{8*}}\right]_8 P_8^*\right)M_8 - \left[\dot{f_{8*}}\right]_8 P_8^* W_{,8}$$

From (30), we arrive at the equations

$$\frac{dW_8(t)}{dt} = \int_0^\infty ([h_*]_8\beta_8(a_8) - \mu_8(a_8))m_8(t,a_8)da_8$$

$$-\left([f_{8*}] + [f_{8*}]_8 + [f_{8*}']_8 P_8^*\right) M_8 + \left(\frac{[h_*']_8}{[h_*]^8}\, P_8^*(0) - f_{8*}' P_8^*\right) W_8,$$

(31)

$$\frac{dW_8(t)}{dt}\left[g'_*\right]_8 Q_8^*\, M_{.8}$$

By the mean value theorem,

$$\int_0^\infty ([h_{"*}]_8 \beta_8(a_8) - \mu_8(a_8)) m_8(t, a_8) da_8 = ([h_*]_8 \beta_8(\theta_8) - \mu_8(\theta_8)) M_{,8}$$

Where $0< \theta_8 = \theta_8(t) < \infty$.

And if a $^8_8$[a $^8_8$]2<4 [a $_8$]2<4 a $_8$e,

Re $\lambda_8[\lambda_8]_8 = \text{Re}\lambda_8[\lambda_8]_8 = -\frac{8}{8} a_8[\text{a2}\ ]_8 < 0$.

Let ⊿=max$_{8<u<\infty}$ |h $_*$ [h $_*$] $_8$β $_8$($a_8$) − $\mu_8(a_8)$|. Then|$n_8$(t)| ⊿∀t . Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium(ρ $^*$ (a), $\boldsymbol{Q}^*$) is asymptotically stable if either a $^8_8$ <4 a $_8$ e and [a $_8$] > 1⊿,or [a $^8_8$] [a $_8$]e and (a $_8$ − $\sqrt{(\text{a}^8_8)_8 - 4(\text{a}_8)_8\,\text{e}}$ > 1⊿

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{\text{a}}{\text{a}+8}\right)$ and β $_8$($a_8$) = ae-$^{ma}$ Then ⊿=[$\mu_0$] + 2 Simple calculation leads to

$$\int_0^{\infty} \beta_8(a_8)\mathrm{e}^{-\int_0^{\mathrm{a}} \mu_8(\tau_8) + d\tau_8 \mathrm{da}} = \frac{1}{(\mathrm{m}+\mu_0+1)^8} + \frac{1}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_8 + [\mu_0]$ 0.5. Assume$[f_8](x_8) = K_8 x_8$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_8(x_8) = [0x_8]/(2 + x_8)$; the functional responses $[[f_8](x) = [0]\, x_8/(1 + x_8)$and $h_8(x_8)$=$(1+\omega_8 x_8)/(1 + x_8)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_8 + [\mu_0] + [\mathrm{K}][P_0])}{[m_8]_8 + (\mu_0)_8 + [\mathrm{K}][P_0])_8^3\ [\mathrm{K}]_8 [P_0]_8^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_8$i e $[P_0] < b/(\delta_8 - b_8)$then the equilibrium$(P_0)$ $[P_0]_8(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_8$

$[f_{8*}] = K_8\ \mathrm{P}^*\ [f_{8*}] = \frac{\sigma \mathrm{Q}_8^*}{8+\mathrm{Q}_8^*} \quad [h_*] = \frac{8+\omega_8 \mathrm{Q}_8^*}{8+\mathrm{Q}_8^*}$

$[f_{8*}] = K_8\ [\mathrm{f}_{8*}^8] = \frac{\sigma_8}{(8+\mathrm{Q}_8^*)^8}[\ \mathrm{g}_*^8] = \frac{\sigma_8}{(8+\mathrm{P}_8^*)^8} \quad [\mathrm{h}_*^8] = \frac{\omega_8 - 8}{(8+\mathrm{Q}_8^*)^8}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta$, and $\omega$ need to satisfy

$$\frac{\omega_8}{(m + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 + b_8} + \sigma)^8} + \frac{1\omega_8}{([m_8]_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 + b_8} + \sigma_8)^3} < 1$$

And

$$\frac{1}{(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^8} + \frac{1\omega_8}{(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8} + (m_8 + [\mu_0]_8 + \frac{K_8 b_8}{\delta_8 - b_8} + \sigma_8)\, Q_8^*)^3$

$= m_8 + [\mu_0]_8 + 3 + \frac{K_8 b_8}{\delta_8 - b_8}(1 + Q^*)(1 + \omega_8 Q_8^*$

Now

$$\left[a_1^1\right] - 4\left[a_8\right]e = \left(\frac{1K_8 b_8}{\delta_8 - b_8}\right)^8 + \frac{4K_8 b_8 \sigma_8 Q_8^*}{(\delta_8 - b_8)(8 + Q_8^*)} + \frac{\sigma_8^8 Q_8^{*8}}{(1 + Q_8^*)^8}$$

$$\frac{4\, b_8 \sigma_8 (\delta_8 - b) Q^*}{\delta_8 (1 + Q_8^*)^8} - \frac{4\, (\delta_8 - b_8)^8 (1 - \omega_8) Q_8^*}{\delta_8 (1 + \omega_8 Q_8^*)(1 + Q_8^*)} P^*(0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_8\, ae_8^{mu}\ [\mu_0]_8 - \frac{a_8}{a_8 + 8}$$

Hence if $[a_8^8 > 4(a_8)_8 e$ and

$$[a_8] = \frac{1K_8 b_8}{\delta_8 - b_8} + \frac{\sigma_8 Q_8^*}{1 + Q_8^*} > 1\Delta$$

**The nontrivial equilibrium is stable.**

IF we define

$$[a_8]_8 = [f_{8*}]_8 + [f_{8*}]_8 + [f_{8*}] P_8^*$$

$$[a_8]_8 = \left[f_{8*}^{8}\right] P_8^* - \frac{h_*^8}{h_*} P^*(0)$$

$$e = [g_*^8]_8 Q_8^*$$

$$\eta_8(t) = \left[h_*\right]_8 \beta_8 \left(\theta_8(t)\right) - \mu_8(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\left[a_8\right] M_8 - [a_8]_8 W_8 - [n(t)] M_8$$

$$w_8(t) = eM_8$$

Write

$$A = \begin{pmatrix} -[a_8]_8 - [a_8]_8 \\ e \qquad 0 \end{pmatrix}$$

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| $B=(t)=\begin{pmatrix} [\eta(t)]_8 & 0 \\ 0 & 0 \end{pmatrix}$ |
| And |
| $X(t)=\begin{pmatrix} M_8 \\ W_8 \end{pmatrix}$ |
| The solution of (32)is |
| $X_8(t) = e^{A,t}[X]_8(0) + \int_0 e^{At}$ |
| This equation notwithstanding its importance in all the models stand out for its importance for the following models<br><br>to which it is applied: |
| **Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016** |

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| (376)These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field. |
| (377)The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets. |
| (378) Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state. |
| (379) Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands. |
| (380)The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI |
| (381) The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9]. |
| (382)This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive. |

(383)More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

(384)Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.

(385)Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.

(386) The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.

(387)This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.

(388)The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.

(389)This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.

(390) For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.

(391)The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(392)Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.

(393)We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.

(394)When this is true—as is often the case—the system serves as (=) a "bath" for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called **thermalization.**

(395)However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.

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| (396)Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1). |
| (397)Bringing them to thermal equilibrium requires (e) an external bath. |
| (398) But so far, there has been no rigorous proof of (e) the existence of **such many-body localization**. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption. |
| (399)The concept of **localization** was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations). |
| (400)The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century. |
| (401)But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information. |
| (402)Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment. |
| (403)For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood. |
| (404) It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7]. |
| (405) Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature. |
| (406) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators. |
| (407)However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8]. |
| (408)Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems. |
| (409) Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL. |
| (410)As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked. |
| (411)His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice. |
| (412) An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site. |

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| (413)This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself). |
| (414)This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition. |
| (415)However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime: |
| (416) These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects. |
| (417)It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies). |
| (418) This is a realistic scenario for (e) all known systems. |
| (419)On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects. |
| (420)On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other. |
| (421)No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true. |
| (422) His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found. |
| (423)Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization? |
| (424)Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects? |
| (425)Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions. |
| (426)So in those systems, the quantum phase transition does indeed exist. |
| (427)But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension. |
| (428)Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters. |

# PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION ATTRITION MODELS

## MODULE NINE

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Model is applied to following formulations Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society**

**Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice ($\alpha$=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP This paper is divided into four parts. In part A, some general considerations about (e&eb) gravitational radiation are followed by (e) a treatment of the scalar wave equation in the manner later to be applied to (e&eb) Einstein's field equations. In part B, a co-ordinate system is specified which is suitable for (e) investigation of outgoing gravitational waves from (e) an isolated axi-symmetric reflexion-symmetric system. The metric is expanded in (eb) negative powers of (e) a suitably defined radial co-ordinate r, and the vacuum field equations are investigated in detail. It is shown that the flow of information to infinity is controlled by (e&eb) a single function of (e) two variables called the news function. Together with initial conditions specified on a light cone, this function fully defines (eb) the behaviour of the system. No constraints of (e) any kind are encountered. In part C, the transformations leaving (eb) the metric in the chosen form are determined. An investigation of the corresponding transformations in (eb) Minkowski space suggests (eb) that no generality is lost by (e) assuming that the transformations, like the metric, may be expanded in (eb) negative powers of r. Gravitational Waves in General Relativity VII Waves from Axi-Symmetric Isolated Systems Royal Society Proceedings H. Bondi, M. G. J. van der Burg, A. W. K. Metzner Published 21 August 1962.DOI: 10.1098/rspa.1962.0161 STOP There is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant $\Lambda$, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive $\Lambda$ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the $\Lambda>0$ generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible. Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of (e) GW signals from BBHs with component spins $\chi_{1,2}\in[-0.99,0.99]$ aligned (e&eb) with the orbital angular momentum, component masses $m_{1,2}\in[2,48]M_\odot$, and (e&eb) total mass $M_{total}\leq 50\ M_\odot$. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin**

**bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins >+0.6 and component masses ≳20 M⊙. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201**

## 54. MODEL FORMULATION

Let $p_9(t, a_9)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(g_9)Q_9$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_9(0) = 0, \quad g_9(P_9) > 0, \ P_9 \ 0, \qquad \lim_{P_9 \to \infty} g_9(P_9) = (K_9)_9 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_9 + (\rho_a)_9 + D_9(a_9, P_9, Q_9)\rho_9 + 0, \qquad 0 < a_9 \ \infty,$$

$$\rho_9(t,0) = \int_0^\infty n_9(a_9, P_9, Q_9)\rho_9(t_9, a_9)da_9$$

Where $D_9$ is the mortality function and $n_9$ is the fertility function.

Put $_9$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_9}{dt}\right]_9 = x_9\left[\frac{dQ_9}{dt}\right]_9 - x_9\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_9}{dt}\right]_9 = x_3\left[\frac{dQ_3}{dt}\right]_9 - x_4\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_9}{dt}\right]_9 - x_6\left[\frac{dQ_3}{dt}\right]_9$$

Where $x_9, x_9, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_9$ and $g_9(P_9)Q_9$

Now when you add the term $(b_9Q_9)$ with $g_9(P_9)(Q_9)$ we are infect see is that $(b_9Q_9)$ is credited incrementally part of $g_9(P_9)(Q_9)$ or instant nearly, when subtracted from $\frac{dQ_9}{dt}$ , it means that $g_9(P_9)(Q_9)$ is slowly eaten by $\frac{dQ_9}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_9 = \mu_9(a_9) + (f_9)_9(P_9)_9 + (f_9)_9(Q)_9 \quad (2)$$

$$n_9 = \beta_9(a_9)h_9(Q_9), \quad (3)$$

Where $\mu_9$the age- dependent debit rate whit is $\mu_9(a_9)$ 0 $\forall a > 0$; $(f_9)_9$ is the density- dependent death rate with $f_9(0) = 0$, $(f_9^9)_9(P_9)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_9)_9(P_9) = (k_9)_9$ ; $f_9$ is the functional response of creditors to prey with $[f_9]_9(0) = 0, (f_9^9)_9(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_9)_9(Q_9) = [k_9]$ ; $\beta$ is the age dependent birth

rate with $\beta(a) \geq 0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h'(Q) \leq 0, \forall Q \geq 0$, and $\lim_{Q\to\infty} h(Q) = k_3 \geq 0$.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_9 + (\rho_a)_9 = -\mu_9(a_9) + (f_9)_9 P_9) + \big(f_9)_9(Q_9)\big)(\rho_9)_9 \tag{4}$$

$$\rho_9(t_9, 0) = h_9(Q_9) \int_0^\infty \beta_9(a_9)\rho_9(t, a_9)da, \tag{5}$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(P_9)Q_9 \tag{6}$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_9$ and $Q_9$ of (4-6) satisfy

$$0 \leq \rho_9(t, a_9) \leq \rho_9^*(t, a_9),$$

$$0 \leq Q_9(t) \leq Q_9^*(t),$$

Where $\rho_9^*(t, a_9)$ and $Q_9^*(t)$, are solutions of

$$(\rho_t^*)_9 + (\rho_a^*)_9 = -(\mu_9(a_9) + (f_9)_9(P_9^*) + (f_9)_9 Q_9^*))\, \rho_9^*, \tag{7}$$

$$\rho_9^*(t, 0) = \int_0^\infty \beta_9(a_9)\, \rho_9^*(t, a_9)da_9 \tag{8}$$

$$\frac{dQ_9^*}{dt} = -b_9 Q_9^* + g_9(P_9^*)Q_9^*, \tag{9}$$

With $P_9^*(t) = \int_0^\infty \rho_9^*\ (t, a_9)da_9$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_9^*(t)}{dt} = P_9^*(t)\big(r_9 - (f_9)(P_9^*) - (f_9)_9(Q^*)\big), \quad (10)$$

$$\frac{dQ_9^*(t)}{dt} = Q_9^*(t)(-b_9 + g_9(P_9^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_9(a)e^{-\int_0^a \mu_9(\tau)dt}e^{-ra}\ da_9 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_9^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(t)dt}\ da_9 > 1, \quad (H1)$$

Which guarantees $r_9 > 0$.

*Remark.* If $\beta_9(a) = [\beta_0]_9 a_9 e^{-\alpha_9 a_9}, \mu_9(a_9) \equiv \mu_9$ then the solution of (12) is

$$r_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_9]_9(P_9) = 0, [f_9](Q_9) = \lambda_9 Q_9\ h_9(Q_9) = 1, \quad and\ \ g_9(P_9) = c_9 P_9$$

$$[\rho_9]_9 + [\rho_a]_9 = -(\mu_9 + \lambda_9 Q_9)\rho_9$$

$$\rho_9(t,0) = \int_0^{\infty} [\beta_0]_9 a_9 e^{-\alpha_9 a_9} \rho_9(t,a_9)\, da_9$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + c_9 P_9 Q_9$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_9}{dt} = \omega_9 P_9 - \lambda_9 P_9 Q_9$$

$$\frac{dQ_9}{dt} = b_9 Q_9 - c_9 P_9 Q_9$$

Where $\omega_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$

THE EQUILIBRIUM $(\rho_0)_9(a_9), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_9(a_9), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_9 + (\rho_a)_9 = -\big(\mu_9(a_9) + (f_9)_9(P_9)\big)\rho,$$

$$\rho_9(t,0) = \int_0^{\infty} \beta_9\,(a_9)\rho_9(t,a_9)da_9 \qquad (13)$$

The solution of (13) has the form of

$$\rho_9(t,a_9) = Ce_9^{r(t-a_9)} - \int_0^a (\tau)d\tau - \int_0^9 (f_9)_9\big(P_9(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_9(t) = C_9 e^{r_9 t - \int_0^t (f_9)_9(P_9(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_9^{-9}(r)$ if $k_9 > r$. This is equivalent to

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{(k_9)_9 a_9} da_9 < 1.$$

In fact, if we define

$$R_9(s_9) = \int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau)d\tau} e^{-s_9 a_9}\, da_9$$

$R(s)$ is a decreasing function. $R_9(r_9) = 1$ And $R_9(k_9)_9 < 1$ imply $r_9 < (k_9)_9$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_9(a_9), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{k_9 a_9} da_9 \quad < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_9 \equiv \int_0^\infty (\rho_0)_{99}(a_9)da_9$ *satisfies*

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{-(f_9)_9(P_0)a_9} da \quad = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_9(a_9), 0)$

Let $u_9(t,a_9) = \rho(t,a_9) - \rho_0(a_9)$ Then

$$N_9(t) \equiv \int_0^\infty u_9\,(t, a_9)da_9 = \int_0^\infty \rho_9(t, a_9)da_9 - \int_0^\infty (\rho_0)_9(a_9)da = P_9(t) - P_0)_9$$

And

$$(u_9)_9 + (u_a)_9 = (\rho_t)_9 + (\rho_a)_9 - (\rho_0^9)_9(a_9)$$

$$= -\mu_9(a_9) + (f_9)_9(P_9) + (f_9)_9(Q_9))\,(u_9 + (\rho_0)_9) + \mu_9(a_9)(\rho_0)_9 + (f_9)_9(P_0)_9(\rho_0)_9$$

$$= -\mu_9(a_9)u_9 - (f_9)_9(P_9) + (f_9)_9(Q_9))\,u_9 - (\rho_0)_9\,(f_9)_9(P_9) - (f_9)_9(P_0)_9$$

$$+ (f_9)_9(Q_9)), \qquad (15)$$

Where $(\rho_0^9)_9(a_9)$ denotes $d(\rho_0)_9(a_9)/da_9$

Expansion $(f_9)_9(P_9)$ and $(f_9)_9(Q_9)$ at $(P_0)_9$ and 0 yield respectively

$$(f_9)_9(P_9) = (f_9)_9((P_0)_9 + N_9)$$

$$= (f_9)_9(P_0)_9 + (f_9^9)_9(P_0)_9N_9 + o_9(N_9), \qquad (16)$$

$$(f_9)_9(Q_9)$$

$$= (f_9^9)_9(0)Q + o_9(Q_9) \qquad (17)$$

Write $f_9(P_0) = f_{90}, f_9'(P_0) = f_{90}'$, and $f_9'(0) = f_{90}'$ and substitute (16) and (17) into (15) to obtain

$$(u_9)_9 + (u_a)_9 \approx -(\mu_9(a_9) + (f_{90})_9)u - (f_{90}^9)_9(\rho_0)_9(N_9) - (f_{90}^9)_9(\rho_0)_9(Q_9)$$

$$u_9(t, 0) = \rho_9(t, 0)(\rho_0)_9(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t, a) + (\rho_0)_9(a_9)da_9 - (\rho_0)_9\,(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t, a_9)da_9 + (h_9(Q_9) - 1)(\rho_0)_9\,(0) \qquad (18)$$

Expanding

$$h_9(Q_9) = 1 + (h^9)_9(0)Q_9 + 0(Q_9) \equiv 1 + h_0'Q_9 + 0Q_9$$

And substituting it into (18), we arrive at

$$u_9(t,0) \approx \int_0^\infty \beta_9(a_9)u_9(t,a_9)da_9 + (h_{a_9}')(\rho_0)_9Q_9$$



PREDATOR-PREY POPULATION MODELS

The linearization is the following :

$$[u_9] + [u_9] = -(\mu_9(a_9) + [t_{90}]_9u_9 - [\alpha_9]_9(a_9)N_9 - [\alpha_9]_9(a_9)Q_9$$

$$u_9(t,0) = \int_0^\infty \beta_9(a_9)u_9(t,a_9)da_9 + [\xi_9]Q_9 \qquad (19)$$

$$\frac{dQ_9}{dt} = -b_9Q_9$$

where

$$[\alpha_9]_9(a_9) = [t_9^9a]_9[p,o]_9(a_9)d2(a_9) = [f_{90}^9]_9(\rho\, o\,)_9(a_9)\xi = [\mathrm{h\ o}^9]_9[\rho\mathrm{o}]_9(\mathrm{o}) \text{ and } \mathrm{b}_9 = \mathrm{b}_9 - g_9(P_9)$$

If $g_9(P_0) > b$, clearly, $(\rho_0)_9(a_9), 0)$ is unstable and the predator population is persistent. Assume $g_9(P_0)_9 < b$.

Then, $Q_9(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_9 + u_a = -(\mu_9(a_9) + (f_{90})_9)u_9 - (\alpha_9)_9(a_9)N_9$$

$$= -(\mu_9(a_9) + f_{90})u_9 - \alpha_9(a_9)\int_0^\infty u_9(t,a_9)da_9$$

$$u_9(t,0) = \int_0^\infty \beta_9\,(a_9)u_9(t,a_9)da_9$$

Using the method developed by Cushing [2, 3], we suppose

$$u_9(t,a_9) = \tilde{u}_9(a_9)^{e^{\gamma(t-a_9)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_9}\,da_9$. Then,

A solution of the first differential equation for $\tilde{u}(a_9)$ gives,

$$\tilde{u}_9(a_9) = e^{-\int_0^a \mu_9(\tau_9)d\tau-(f_{90})_9 a}\left(\tilde{u}_9(0) - w\int_0^a e^{\int_0^\tau \mu_9(s_9)ds+(f_{90})}\tau_9\,(\alpha_9)_9(\tau_9)e^{\gamma_9\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\right)\tilde{u}_9(0)$$

$$+w\int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\int_0^a (\alpha_9)_9(\tau_9)e^{-\gamma_9\tau_9}e^{(f_{90})_9\tau_9+\int_0^\tau \mu_9(s_9)ds}d\tau_9\,da_9$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}da\right)(\tilde{u}_9(0))_9$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\int_0^a e^{(f_{90})_9\tau_9+\int_0^\tau \mu_9(s_9)ds+y\tau}\alpha_9(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(t_9)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_9\,(a_9)Y_9(\gamma_9,a_9)da_9\right)u_9(0)$$

$$+w\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)\int_0^a(\alpha_9)_9(\tau)_9(Y^{-9})_9(\gamma_9,\tau_9)d\tau_9\,da_9=0,$$

(24)

$$\left(\int_0^{\infty}Y_9(\gamma_9,a_9)da_9\right)\tilde{u}_9(0)$$

$$+\left(-1-\int_0^{\infty}Y_9(\gamma_9,a_9)\int_0^a(\alpha_9)_9(\tau)_9(Y^{-9})_9(\gamma_9,\tau_9)d\tau_9\,da_9\right)w=0.$$

Equation (20) has a solution of the form of $u_9(t,a_9)=u_9(a_9)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_9(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_9(\gamma_9,a_9)=\int_0^a(\rho_0)_9(\tau_9)Y^{-9}\,(\gamma_9,a_9)d\tau_9$. We obtain the characteristic equation

$$f_{90_9}^{l}\left(\int_0^{\infty}Y_9(\gamma_9,a_9)da_9\right)\left(\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)I_9(\gamma_9,a_9)da_9\right)$$

$$+\left(1-\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)da\right)\left(1+(f_{90}^{l})_9\int_0^{\infty}Y_9(\gamma_9,a_9)I_9(\gamma_9,a_9)da_9\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_9(a_9) = (\rho_0)_9(0)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - (f_{90})_9 a_9}$$

$$I_9(\gamma_9, a_9) = \int_0^a (\rho_0)_9(\tau_9)e^{\gamma_9\tau_9 + (f_{90})_9\tau + \int_0^\tau \mu_9(s_9)ds_9}d\tau$$

$$= \frac{(\rho_0)_9(0)}{\gamma_9}(e^{\gamma_9 a_9} - 1)$$

Let

$$E_9(a_9) = e^{-(f_{90})_9 a_9 - \int_0^a \mu_9(\tau_9)d\tau_9}$$

Then $[Y_9]_9[\gamma_9]_9, a_9) = E_9(a_9)e^{-\gamma_9 a_9}$. by a simple calculation,

$$\int_0^\infty \beta_9(a_9)Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma}\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}\ da\right);$$

$$\int_0^\infty Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma_9}\left(\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}\ da_9\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_9(0)}{\gamma_9}[f_{90}^9]\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}da_9\ (\int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9)$$

$$+\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right).$$

$$\left(1+\frac{(\rho_0)_9(0)}{\gamma_9}(f_{90}^{'})_9\left(\int_0^{\infty}E_9(a_9)da_9-\int_0^{\infty}E_9(a_9)e^{-\gamma_9 a_9}da_9\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right)(1+\frac{(f_{90}^{'})_9}{\gamma_9}(\rho_0)_9(0)\int_0^{\infty}E_9(a_9)da_9)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (94), to

$$1-\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)da_9=1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\neq 0,$$

Since $\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)\int_0^{a}\alpha_9(\tau_9)Y^{-9}(\gamma_9,\tau_9)d\tau_9 da_9>0$. Therefore, the characteristic solution is

$$\gamma_9=-(f_{90}^{'})_9(\rho_0)_9\int_0^{\infty}E_9(a_9)da_9.$$

Since $f_9^{'}(P_0)_9>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_9),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_9),0)$ which is related to the extinction of predators (debit summation types: that which acts as a resource)and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_9 + (\mu_9(a_9) + [f_9]_9(P_9^*) + [f_9]_9(Q_9^*))P_9^* = 0, \quad (26)$$

$$p_9^*(0) = h_9(Q_9^*)\int_0^\infty \beta_9(a_9)p_9^*(a_J)p_9^*(a_9)da_9 \quad (27)$$

$$g_9(P_9^*) - b_9 = 0, \quad (28)$$

Where $P_9^* = \int_0^\infty p_9^*(a_9)\, da_9$

Equation (28) yields $P^* = g^{-9}(b)$,

From (26),

$$p_9^*(a) = p_9^*(0)e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}$$

Substituting into (27), we need to solve

$$\left(1 - h_9(Q_9^*)\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}da_9\right)p_9^*(0) = 0$$

For $Q_9^*$

Fix $p_9^*(0)$ and $P_9^*$ Let

$$H_9(z_9) = 1 - h_9(z_9)\int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9}\, da_{,9}$$

And assume that

$$[k_3]_9 \int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(g_9^{-9}(b_9))a_9 - [k_9]_9 a_9}\, da_9 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_9(z_9)] > 0.$$

Clearly,

$$\frac{dH_9}{dz_9} = h_9(z)\left([f_9']_9(z_9)\right)\int_0^\infty a_9\beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - f_9(P_9^*)a_9 - f_9(z_9)a_9} da_{\square}$$

$-h'(z_9)\int_0^\infty \beta_9(a_9)\, e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9}\ da_9$

y the assumptions on $\left[f_9^9\right]_9 and\ [h^9]^9 dH_9/dz > 0$. Hence, if

$$H_9(0) = 1 - \int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9} da_9 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_9(0) < 0$ follows if $g_9^-(b_9) < [f_9^{-9}]_9(r_9)$, where $r_9$ is the solution of (12).

Indeed, since

$$R_9(s_9) = \int_0^{\infty} \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 e^{-y\frac{a}{9}} da_9}$$

Is a decreasing function of $s_9$ when $P_9^* = [g^{-9}]_9(b_9) < [f_9^{-9}]_9(r_9)$, that is $[f_9]_9(P_9^*) < r$, $R\big(f_9(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_9(g^{-9}(b))) > 1, i.e.,$

$\int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9\left(g_9^{-9}(b)\right)a_9} > 1.$ (H4)

Then the system (4)-(6) has a unique positive equilibrium $(P_9^*(a_9), Q_9^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_9^*(a_9), Q_9^*)$ is more complicated. We again consider perturbations $m_9(t_9 a_9) = p_9(t_9, a_9) - p_9^*(a_9)$ and $W_9(t) = Q_9(t) - Q^*$.Then

$M_9(t) \equiv \int_0^{\infty} m_9(t, a_9) da_9 = P_9(t) - P_9^*,$

$$[m_9]_9 + [m_a]_9 = [p_9]_9 + [p_a]_9 - \left[p^*\right]_9(a_9) \approx -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9$$

$$-[f_{9*}^9]_9 p_9^*(a_9)M - [f_{9*}^9]_9 p_9^*(\mathrm{a})W_9,$$

$$m_9(t,0)p_9(t,0) - p_9^*(0) \approx h_* \int_0^{\infty} \beta_9(a_9)m_9(t,a_9)\, da + \frac{[h_*^9]_9}{[h_*]_9}\, p_9^*(0)W_9$$

$$\frac{dW_9}{dt} = \frac{dQ_9}{dt} \approx [g_*^9]_9 Q_9^* M_{,9}$$

Where $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $[f'_{9*}]_9$ denoting $[f'_9(P^*)]_9 [f'_{9*}]_9$ denoting $[f_9]_9$ $(Q_9^*$, $h_9^*$ denoting $h'_*$ denoting $h'_9$ $(Q_9^*)$, and $[g'_*]_9$

Denoting $[g']_9(P_9^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_9]_9 + [m_a]_9 = -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9 - [f'_{9*}]_9 [p_9^*](a_9)M_9 - [f'_{9*}]_9 p_9^*(a_9)W_9$$

$$m_9(t,0) = [h_*]_9 \int_0^\infty \beta_9(a_9)m_9(t,a_9)da_9 + \frac{[h_*^9]_9}{[h^*]_9}[p^*]_9(0)\, W_{,9} \qquad (30)$$

$$\frac{dW_9}{dt} = [g_*^9]_9 [Q_9^*]M_9$$

Integrating $[m_9]_9(t,a_9)$ with respect to a from 0 to ∞, we have

$$\frac{dM_9(t)}{dt} = \int_0^\infty [mt]_9(t,a_9)da_9$$

$$= -\int_0^\infty [m_a]_9\,(t,a_9)da_9 - \int_0^\infty \mu(a_9)m(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + [f'_{9*}]_9 \int_0^\infty p_9^*(a_9)da_9\right)M_9 - [f_{9*}^9]_9 \int_0^\infty p_9^*(a_9)da_9 W_9$$

$$= m_9(t,0) - \int_0^\infty \mu_9(a_9)\, m_9(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + [f'_{9*}]_9 P_9^*\right)M_9 - [f'_{9*}]_9 P_9^* W_{.9}$$

From (30), we arrive at the equations

$$\frac{dW_9(t)}{dt} = \int_0^{\infty} ([h_*]_9\beta_9(a_9) - \mu_9(a_9))m_9(t,a_9)da_9$$

$$-\left([f_{9*}] + [f_{9*}]_9 + \left[f_{9*}^{'}\right]_9 P_9^*\right) M_9 + \left(\frac{\left[h_*^{'}\right]_9}{[h_*]^9}\ P_9^*(0) - f_{9*}^{'} P_9^*\right) W_9,$$

(31)

$$\frac{dW_9(t)}{dt}\left[g^{'}{}_*\right]_9 Q_9^*\, M_{.9}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{"*}]_9\beta_9(a_9) - \mu_9(a_9))m_9(t,a_9)da_9 = ([h_*]_9\beta_9(\theta_9) - \mu_9(\theta_9))M_{,9}$$

Where $0< \ \theta_9 = \theta_9(t) < \infty$.

And if $a_9^9[a_9^9]2<4\ [a_9]2<4\ a_9e$,

$\text{Re}\,\lambda_9[\lambda_9]_9 = \text{Re}\lambda_9[\lambda_9]_9 = -\frac{9}{9}a_9[a2\ ]_9 < 0.$

Let $\Delta=\max_{9<u<\infty}\left|h_*[h_*]_9\beta_9(a_9) - \mu_9(a_9)\right|$. Then$|n_9(t)|$ ⍰⊿∀t ⍰0. Hence we obtain the following result.

**THEOREM 4**

Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_9^9 <4\ a_9e$ and $[a_9] > 1\Delta$,or $[a_9^9]$ ⍰4$[a_9]e$ and $(a_9 - \sqrt{(a_9^9)_9 - 4(a_9)_9e) > 1\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+9}\right)$ and $\beta_9(a_9) = ae^{-ma}$ Then $\Delta=[\mu_0] + 2$ Simple calculation leads to

$$\int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9) + d\tau_9 da} = \frac{1}{(m + \mu_0 + 1)^9} + \frac{1}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_9 + [\mu_0]$ 0.5. Assume$[f_9](x_9)= K_9 x_9$ Then the hypothesis (H9) is satisfied We take the conversion rate function $g_9(x_9)= [0x_9]/(2 + x_9)$; the functional responses $[[f_9](x) = [0]\, x_9/(1 + x_9)$and $h_9(x_9)=(1+\omega_9 x_9)/(1 + x_9)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_9 + [\mu_0] + [K][P_0])}{[m_9]_9 + (\mu_0)_9 + [K][P_0])_9^3\ [K]_9[P_0]_9^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_9$i e $[P_0] < b/(\delta_9 - b_9)$then the equilibrium$(P_0)$ $[P_0]_9(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_9$

$[f_{9*}] = K_9\ P^*\ [f_{9*}] = \frac{\sigma Q_9^*}{9+Q_9^*}$ $[h_*] = \frac{9+\omega_9 Q_9^*}{9+Q_9^*}$

$[f_{9*}] = K_9\ [f_{9*}^9] = \frac{\sigma_9}{(9+Q_9^*)^9}\ [g_*^9] = \frac{\sigma_9}{(9+P_9^*)^9}$ $[h_*^9] = \frac{\omega_9 - 9}{(9+Q_9^*)^9}$

In order to have (H3) and (H4) hold,$m, \mu_0, K, b, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_9}{(m+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9+b_9}+\sigma)^9}+\frac{1\omega_9}{([m_9]_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9+b_9}+\sigma_9)^3}<1$$

And

$$\frac{1}{(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9})^9}+\frac{1\omega_9}{(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9}+(m_9+[\mu_0]_9+\frac{K_9 b_9}{\delta_9-b_9}+\sigma_9)\,Q_9^*)^3$$

$$=m_9+[\mu_0]_9+3+\frac{K_9 b_9}{\delta_9-b_9}(1+Q^*)(1+\omega_9 Q_9^*$$

Now

$$\begin{bmatrix}1\\1\end{bmatrix}-4\begin{bmatrix}a_9\end{bmatrix}e=\left(\frac{1K_9 b_9}{\delta_9-b_9}\right)^9+\frac{4K_9 b_9\sigma_9 Q_9^*}{(\delta_9-b_9)(9+Q_9^*)}+\frac{\sigma_9^9 Q_9^{*9}}{(1+Q_9^*)^9}$$

$$\frac{4\,b_9\sigma_9(\delta_9-b)Q^*}{\delta_9(1+Q_9^*)^9}-\frac{4\,(\delta_9-b_9)^9(1-\omega_9)Q_9^*}{\delta_9(1+\omega_9 Q_9^*)(1+Q_9^*)}P^*(0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_9 \, ae^{mu}_9 \quad [\mu_0]_9 - \frac{a_9}{a_9+9}$$

Hence if $[a^9_9 > 4(a_9)_9 e$ and

$$[a_9] = \frac{1K_9 b_9}{\delta_9 - b_9} + \frac{\sigma_9 Q^*_9}{1 + Q^*_9} > 1\Delta$$

Some of the examples to which the model has been applied are:

**New look at the galaxy-shaping power of black holes**

(596) "It creates (eb) these bubbles, and the bubbles keep the plasma hot.

(597) That's what prevents galaxies from (e) becoming even bigger than they are now."

(598) Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.

(599) In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.

(600) Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.

(601) "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.

(602) "It's like nature's thermostat, which keeps these galaxies from (e) growing.

(603) If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."

**A new look at the galaxy-shaping power of black holes**

(604) An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the

(605) McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.

(606) Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.

**Doomed satellite missions**

(607) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.

(608) At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.

(609) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.

(610) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.

(611) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.

(612) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people

(613) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."

(614) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.

(615) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(616) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.

(617) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.

(618) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.

(619) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.

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| (620) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well. |
| (621) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma. |
| (622) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6. |
| (623) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus. |
| (624) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (625) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (626) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (627) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (628) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (629) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (630) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (631) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (632) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (633) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (634) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (635) We know dark matter exists only from (e) its gravitational effects. |
| (636) It holds together galaxies and (e&eb) galaxy clusters. |
| (637) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (638) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |

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| (639) | They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (640) | X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (641) | They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (642) | This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (643) | An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |
| (644) | Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (645) | "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (646) | The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (647) | Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (648) | Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (649) | "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (650) | Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (651) | "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (652) | Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) |

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| consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories. |
| (653) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen. |
| (654) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer. |
| (655) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants. |
| (656) These were treated as (=) "aware" and "unaware" trials, respectively. |
| (657) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it. |
| (658) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another. |
| (659) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion. |
| (660) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness. |
| (661) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said. |
| (662) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness." |
| (663) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner. |
| (664) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain. |
| (665) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said. |
| (666) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

(667) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter.

(668) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined.

(669) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions.

(670) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.

(671) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.

(672) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.

(673) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.

(674) Noda and his colleagues instead studied the change in emissions over (e&eb) time.

(675) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.

(676) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).

(677) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.

(678) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.

(679) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson’s problem**

(680) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking “palm tree” organization.

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| (681) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society |
| (682) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible? |
| (683) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem. |
| (684) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other. |
| (685) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible. |
| (686) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges. |
| (687) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century. |
| (688) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching. |
| (689) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom. |
| (690) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth. |

**Spheres on trees**

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| (691) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration. |
| (692) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement. |
| (693) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum. |
| (694) By visualizing the problem in this way, the researchers noticed that these particular graphs don’t (e) have lots of branches extending from lots of (e) other branches. |
| (695) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure. |

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| (696) | The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network. |
| (697) | As a result, it provides (eb) an important clue for finding the global minimum. |
| (698) | It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum. |
| (699) | This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large. |

**Thompson's problem 2**

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| (700) | By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society |
| (701) | The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima. |
| **(702)** | The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node. |

**Implications for Thomson's and Smale's problems**

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| (703) | Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before. |
| (704) | "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around." |
| (705) | As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem. |
| (706) | "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view." |
| (707) | In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges. |
| (708) | From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear. |
| (709) | "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK. |
| (710) | "For larger numbers, there will likely be (=) multiple funnels. |
| (711) | Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum." |

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| (712) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem. |
| (713) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said. |
| (714) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org |

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| The nontrivial equilibrium is stable. |
| IF we define $[a_9]_9 = [f_{9*}]_9 + [f_{9*}]_9 + [f_{9*}]\,P^*_9$ |
| $$[a_9]_9 = \left[f^{9}_{9*}\right]P^*_9 - \frac{h^9_*}{h_*}P^*(0)$$ |
| $$e=[g^9_*]_9\,Q^*_9$$ |
| $$\eta_9(t) = \left[h_*\right]_9\beta_9\left(\theta_9(t)\right) - \mu_9\left(\theta(t)\right),$$ |
| The system(32)can be written as the following system: |
| $$[M(t)]\left[a_9\right]M_9 - [a_9]_9W_9 - [n(t)]M_9$$ |
| $$w_9(t) = eM_9$$ |
| Write |
| $$A=\begin{pmatrix}-[a_9]_9 - [a_9]_9 \\ e \qquad\quad 0\end{pmatrix}$$ |

$$B=(t)=\begin{pmatrix} [\eta(t)]_9 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_9 \\ W_9 \end{pmatrix}$$

The solution of (32)is

$$X_9(t) = e^{A,t}[X]_9(0) + \int_0 e^{At}$$

**Note: There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and**

**sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by**

**term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation**

**alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation,**

**appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module**

**LEN NUNNEYS**

## Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

## Module numbered one

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems

and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1965); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be referred to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 6985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_6}{dt} = [N(t)]_6[(F_6)(R(t)]_6 - [M]_6[R(t)]_6$$

$$\frac{d[R(t)]_6}{dt} = [B]_6\,[(R)_6(t-T_6)] - [D]_6[(R(t)]_6 - [N(t)]_6\left[G\left(R(t)\right)_6\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_6}{dt} = a_6[N(t)]_6 - a_6[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_6 = a_3[N(t)]_3 - a_4[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_6 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

irst, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.6b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The

regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 6a is at the

**TIME DELAY**

FIG. 6. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.6$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**DELAY DUE TO PREDATOR MATURATION**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_6}{dt} = N_6(t - T_6)F_6(R_6(t - T_6) - N_6(t)M_6((R_6(t) \qquad (5)$$

$$\frac{dR_6(t)}{dt} = B_6\big(R_6(t) - D_6(R_6(t)\big) - N_6(t)G_6(R_6(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_6}{dR_6} + \frac{N_6 dG_6}{dR_6} - \frac{dB_6}{dR_6} > \frac{N_6 G_6}{2M_6}\left[\frac{dF_6}{dR_6} + \frac{dM_6}{dR_6}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_6/dR_6$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_6$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_6$) and the other derived from the non-delayed or "immediate" processes($I_6$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

ABSOLUTE STABILITY

$$L_6 = \frac{N_6 G_6}{2M_6}\frac{d(6nF_6)}{dR_6} (\geq 0) \qquad (6a)$$

$$I_6 = -\frac{N_6 G_6}{2M_6}\frac{d(6nM_6}{dR_6} (\geq 0) \qquad (6b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_6 > L_6 - I_6$, must be modified ahe $dM_6/dR_6$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_6 > L_6 - I_6, \quad \text{when} \quad P_6^2/4 > I \qquad (8a)$$

$$P_6 > L_6 - I_6 + \frac{(I_6 - P_6^2}{(L_6 - I)} \quad \text{when} \quad P_6^2/4 < I. \quad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

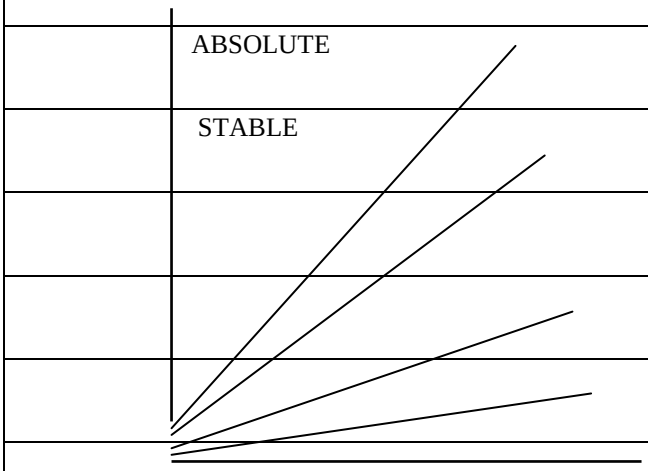



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

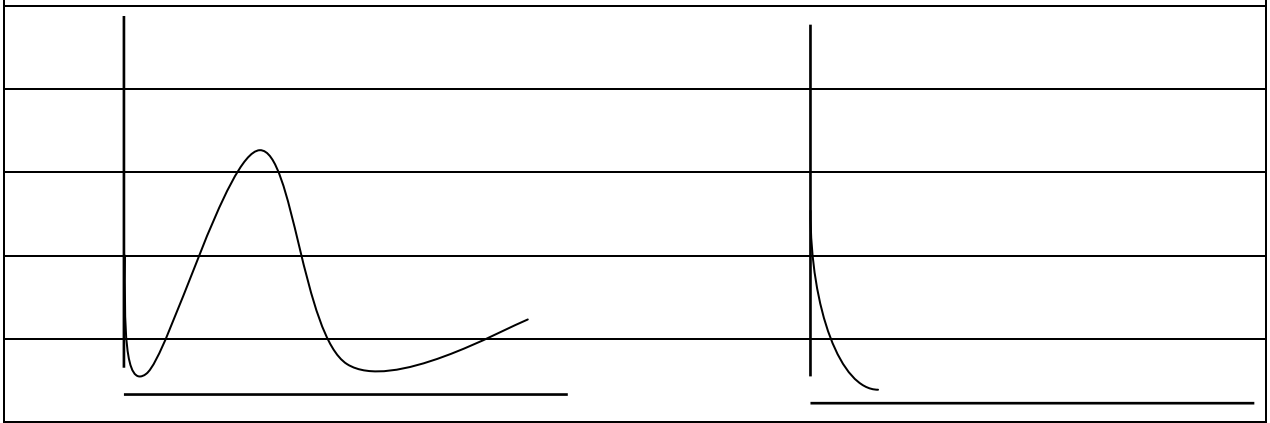
Can never completely offset the lagged effect,$L_6$; condition (8b) ensures that the total debit regulation, $P_6$, must always be positive even when $I_6$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_6$ and $L_6$. As $I_6$ increases relative to $L_6$ the zone of stability is enlarged; however,even when $I_6$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_6$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



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Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (6). Parameters in each case M=6; $dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6 = 0.05$ For (a) $NGdF/dR = 0.3, -N_6G_6dM_6/dR_6 = 0.6$; for (b)$NG_6dM_6/DR_6 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_6t}{dt} = N_6(t - T_6)F_6\big(R_6(t - T_6)\big) - N_6(t)M_6\big(R_6(t)\big)$$

(9)

$$\frac{dR_6t}{dt} = B_6\big(R_6(t - T_6)\big) - D_6\big(R_6(t)\big) - N_6(t)G_6\big(R_6(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) and the effective prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we

could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

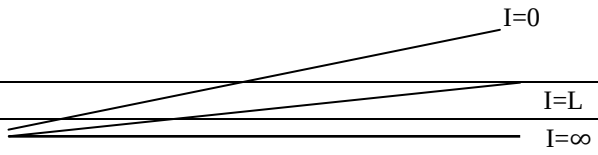
Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_6/dR_6$ is a measure of the control acting through the debit delay and setting $dB_6/dR_6 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_6/dR_6$ is a measure of the control acting through the credit delay although setting $dF_6/dR_6 = 0$ does not negate the effect of a predator delay; the value of $M_6$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_6/dR_6$and normalized by $M_6$ (see6a). The absolute stability of such a system is summarized in Fig. 6 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

TOTAL

ABSOLUTELY STABLE

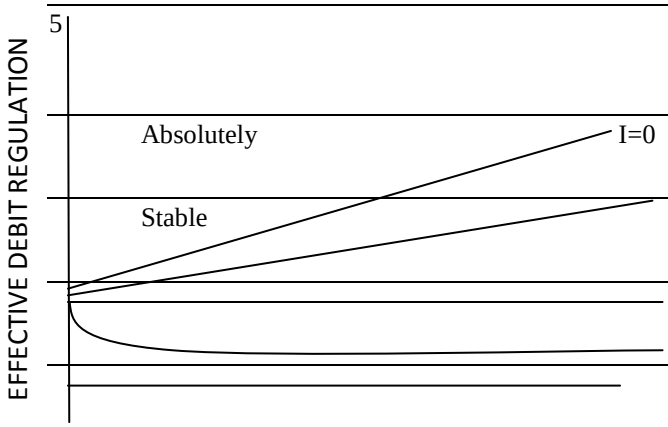



DELAYED DEBIT EFFECT

FIG6. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_6/dR_6 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.



-50 10

DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_6/dR_6 + N_6 dG_6/dR_6 + dB_6/dR_6)/M_6$). The immediate ($I_6$) and lagged ($L_6$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_6$(see 6b), is increased relative to the lagged part, $L_6$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_6 > I_6$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_6$is somewhat large than $I_6$ then the condition constraining prey regulation of recruitment ($dB_6/dR_6 < 0$) is

$$\frac{dD_6}{dR_6} + N_6\frac{dG_6}{dR_6} + \frac{dB_6}{dR_6} > \frac{(L_6 - I_6)M_6}{2} \qquad (10)$$

If $L_6 < I_6$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_6$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_6$

The - $D_6$ positions are defined by letting $\lambda_6 = iZ_6$ and separating (A1.2) into real and imaginary parts:

$$Z_6^2 + V_6 Z_6 \int_0^{\infty} \sin(Z_6 T_6) f(T_6) dT_6 - (X_6 + Y_6) = 0$$

$$Z_6 \left( U_6 - V_6 \int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6 \right) = 0$$

Noting that $X_6 + Y_6 = 0$ is we assume X+Y=0.It follows that $Z_6 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_6 = V_6 \int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6$$

If $U_6 > |V_6|$then given the constraints of $f(T_6)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A6.5)remains

greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

## ABSOLUTE STABILITY WITH A DELAY DUE TO CREDIT CONSOLIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_6^2 + (U_6 - V_6 + M_6)\lambda_6 - M_6\lambda_6 e^{-T_6\lambda_6} + [(U_6 - V_6)M_6 + Y_6] - [((U_6 - V_6 M_6 -]e^{-T_6\lambda_6} = 0$$

(For definitions see Appendix6) We follow the technique outlined in Appendix 6 and let $\lambda_6 = iZ_6$ and separate (A2.6) into real and imaginary parts

$$Z_6^2 + M_6 Z_6 \sin(Z_6 T_6) - (A_6 M_6 + Y_6) + (A_6 M_6 + X_6)\cos(Z_6 T_6) = 0$$

$$(A_6 M_6) Z_6 - M_6 Z_6 \cos(Z_6 T_6) + (A_6 M_6 - X_6)\sin(Z_6 T_6) = 0$$

Where $A_6 = U_6 - V_6$Using (A2.3), substitute for $Z_6^2$ and $Z_6$in (A2.2):

$$W_6^2(X_6^2 - 3M_6 A_6 X_6 - M_6^2(X_6 - Y_6) + 2M_6^3 A_6 + 2A_6^2 M_6^2]$$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_6 = 1 - \cos(Z_6 T_6)$.

In the absence of a delay the system is stable for $A_6 > 0,$ and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_6 > (X_6 - Y_6)/2\ M_6$ Thus a necessary condition for absolute stability is

$$A_6 > [\ X_6 - Y_6/2\ M_6$$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_6$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_6$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_6, A_6) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_6 \le 2$ *is* of interest First consider the extreme values

$$Q_6(0, A_6) = A_6^2\ (X_6 - Y_6)$$

$$Q_6(2, A_6) = (A_6 + 2\ M_6)^2\ (2A_6 M_6 - (X_6 - Y_6))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_6}{\partial A_6} = W_6^2\left(2M_6^3 + 4A_6 M_6^2 - 3M_6 X_6\right) + W_6\left(3M_6 A_6^2 - 2A_6 X_6 + 2M_6\left(2X_6 + Y_6\right)\right)$$

$$+2A_6\left(X_6+Y_6\right)$$

$=2M_6^2W_6^2(M_6+\mathrm{A}_6)+M_6A_6^2W_6+\left(2A_6M_6-X_6+Y_6\right)(A_6W_6+M_6W_6^2$

$+(2-W_6)((2M_6X_6+M_6Y_6)W_6+A_6(X_6+Y_6)),$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_6$quadratic in $W_6$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_6$has no minimum in the range $0 \le W_6 \le 2$.Defining

$Q_6\left(W_6,A_6\right)=a_6W_6^2+c_6 where\ a_6\ b_6\ and\ c_6$ Are functions of $A_6$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_6 \ge b_6 \ge 0$$

And

$$b_6^2 \le 4a_6c_6$$

$$Case(i): Y_6=0$$

$Let\ Y_6 = 0$ And try the substitution A= $X_6 = /2M_6\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_6 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_6 X_6 /$ $2M_6$) is positive for *W<2. Since* $Q_6$ increases with $A_6$(see A2.9), increasing $A_6$will generate no further solutions of $Q_6$=0. Hence the condition A>X/2M defines absolute stability when $Y_6 = 0$.

$Case\ (ii)$: $Y_6 = y_6\ X_6$ ,

$Let\ Y_6 = y_6\ X_6$And try the substitution $A_6 = X_6/(2 + A_6\ y_6)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_6/(2 + n_6 y_6)$ we have

$$W_6^2\ y_6\left[X_6 n_6 \frac{(1 + n_6 y_6)}{(2 + n_6 y_6)} + M_6^2\left(2 - n_6 + n_6 y_6\right)\right]$$

$$-W_6\ X_6\left[\frac{X_6(1 + n_6 y_6)}{M_6^2(2 + n_6 y_6)^2} + 2\ y_6\left(n_6 - 1\right) + \frac{X_6^2(6 + y_6)}{M_6^2(2 + n_6 y_6)} = 0.\ (A2.12)\right.$$

Since $n_6$ and $A_6$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_6$

Defining $E_6 = (2 + n_6 y_6), the\ condition (A2.66?) can\ be\ written\ as$

$[\frac{X_6(E_6 - 6)}{E_6^2}]^2[\frac{X_6^2}{M_6^2} - 4\ y_6\ E_6^2\frac{X_6}{M_6^2} + 4y_6(n_6 - 2)E_6^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in ($X_6 / W_6^2$) can be expressed as

$2E_6[y_6E_6+(y_6E_6(6+y_6)(E_6-n_6))^{6/2}]>\frac{X_6}{M_6^2}>2E_6[y_6E_6-y_6E_6(6+y_6)(E_6-n_6))^{6/2}]$

**ABSOLUTE STABILITY**

Hence if $X_6/M_6^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0$W_6$2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_6-6)}{E_6^2}\left(\frac{X_6}{M_6^2}\right)^2-2y_6\frac{(E_6+n_6^2y_6)}{E_6}\left(\frac{X_6}{M_6^2}\right)-4y_6(E_6-n_6)$$0

And has the solution

$\frac{X_6}{M_6^2}>\frac{E_6}{(E_6-6)}[y_6(E_6+n_6^2y_6)+[y_6^2(E_6+n_6^2y_6)^2+4y_6(E_6-6)(E_6-n_6)]^{6/2}]$

Given $E_6>n_6$this lower limit to $X_6/M_6^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_6/M_6^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_6$in the appropriate range and this limit is further lowered by (A2.64)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_6<n_6(ie.,2A_6M_6^2<X_6-Y_6$)then (A2.6)and

(A2.8)show that $Q_6$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_6}{M_6^2} > 2E_6[y_6E_6 - (y_6E_6(6 + y_6)(E_6 - n_6)]^{6/2}]$$

This equations throws light on the individual general ledger and Nature's general ledger as we call it : see various papers on the issue. powers, which are the energies of Matrka we have already discussed, are not the only obstacles the yogi must overcome. He must, for example, also resist the temptation to rest content with the miraculous yogic powers (siddhi) he acquires in the course of his spiritual develop- ment. Again to do this he must practice Yoga. Similarly, in order to pervade the Fourth State gradually through the other states in the manner proper to practice at the Individual level, the method is the same. He must practise the higher yoga of the Tantras which, turning his mind inwards and freeing it from discursive representations, allows him to penetrate into the Supreme Principle. 312 Once the yogi has attained this contemplative state, his main problem is to make it permanent. In the introverted state the gross external movement of the breath is suspended and with it the activity of the intellect, mind, individualised consciousness, powers of the senses and the ego. 313 When the yogi rises out of this state, he is liable to fall again into the lower order of creation generated by Maya if he does not maintain his awareness of the higher reality he has experienced and allows his awakened, illumined insight to be obscured by the dream-like vision of thought-constructs. 314 Naturally, the yogi must rise out of the introverted condition of suspension. It is inherent in the very nature of reality that it should move out of itself. 315 Pure, universal consciousness initially transforms itself into the vital breath 316 charged with the impression iyasana) of the power of awareness attained through introversion. By attending to the pulse (spanda) of the breath as it moves out of the absolute, the

yogi can develop an intuitive sense of the inherent unity of all he will perceive in the mental and physical spheres created by the outpouring of consciousness. In this way he realises that his own nature is everywhere present in all he perceives and that all things thus reside within him. Blessed with this insight his consciousness remains free and unlimited even at the individual level where the breath, mind, senses and body are active. If the yogi fails to do this, he finds himself once again beset by the strictures of his embodied existence and must, as before, try to pervade all his other states of consciousness with the aesthetic delight (rasa) and wonder (camatkara) of the Fourth State he experienced in contemplation. Again this means that he must strengthen his pure, empowered awareness that his universal nature manifests as all things. 317 In this way he discovers Siva's presence in every sphere of individualised consciousness ranging from the breath to outer objectivity. The yogi's mind then becomes tranquil and undistracted because wherever it may wander, the yogi perceives only Siva, his authentic nature. 318 Consciousness is thus freed of all external referents and the yogi's subjectivity is purified of all identification with the body or anything else that belongs to the objective sphere. The yogi then becomes detached from the opposites of pleasure and pain and is transcendentally free (kevalin). 319 yogis are again, however, liable to fall if he allows himself to get entangled in the play of opposites. This fall is more serious than the others because, although he is caught by the confining restrictions of individual- ised consciousness as before, he is now also affected by karma. Fleeing from pain in the pursuit of pleasure he is bound to act (karma) to minimise one and maximise the other and so is thrown down to the lowest level of embodied subjectivity (sakala). In order to regain his lost state, he must ascend gradually, by Siva's grace, from one order of subjectivity to the next and so free himself progressively of the limitations of the lower levels to gain the greater freedom and expansion of the higher. As he progresses, the objective sphere also evolves from the grossest perceptions of physical objects outside the lowest order to subjectivity, through to the subtler inner, mental

perceptions to finally reach the order of subjectivity that contains objectivity within itself and is free to externalise it at will. 320 The degree to which this process develops depends, as before, on the yogi's awareness of the Fourth State. In consonance with the general principle that the remedy should suit the defect, the yogi is instructed to seek this higher state of consciousness in the wonder (camatkara) or delight (ananda) he feels in moments of intense physical pleasure. At first he experiences this subtle consciousness for an instant in the subjective sphere. If he manages to catch hold of it, it becomes more intense as the cognitive and objective spheres are also gradually pervaded and vitalised by it. Occasions for this practice are, for example, the sense of satisfaction one feels after a good meal or the aesthetic delight one experiences when listening to good music or the pleasure of sexual union with the Tantric consort or even solitary sexual excitation. In these moments of delight the yogi can penetrate momentarily into his own authentic Siva-nature (sambhavavesa) through the empowered contact (saktasparsa) m he makes with it in the freedom of the pure subjectivity of the Fourth State. 322 If the yogi develops his awareness of this higher level of consciousness and maintains it, he eventually experiences it constantly. 323 Clearly, what prevents the yogi from attending to his state of consciousness rather than the circumstances which induce it is the craving for pleasure (abhila§a) born of ignorance — the source of every impurity which clouds consciousness. Craving directs the yogi's attention towards outer, worldly things and so he is caught in the net of thought- constructs. 324 To free himself of his worldly desires and reverse this binding extroversion, the yogi must eradicate its cause. To be freed of all the ups and downs of the path and no longer be tormented by the possibility of a fall, the yogi must see reality perfectly and completely. This insight is itself liberation and the moment it dawns the yogi is instantly freed. This sudden realisation is the goal of Tantric Yoga. E Tantra declares: "He who perceives reality directly, even for the brief moment it takes to blink, is liberated that very instant and never reborn again." 325 Although the yogi's body and mind

continue to function as before, they are like mere outer coverings 326 which contain, but do not obscure, the mighty, universal consciousness which operates through them. The yogi's body is (=) the universe, the senses (=) the energies that (eb, eb+)vitalise it, his mind (=) Mantra, the rhythm of his breath (eb) the pulse of time and his inner nature pure, dynamic consciousness. Raised above all practice, and hence all possibility of falling to lower levels, the yogi realises that he has always been free 327 and that his journey through (e&eb) the **dark land of Maya was nothing but a dream, a construct of his own imagination** The channels of the senses by which the yogi enjoys the pleasure of the objects of sense are those through which he fills the three worlds with the consciousness emitted from his own Heart. 78 **The Doctrine of Vibration T**he senses are outer, subsidiary cycles of conscious energy (anucakra), revolving around the core or primary cycle (mukhyacakra) of pure consciousness. They scintillate with the diversity of sensations and resonate in the field of their awareness like the feathers of a peacock unfurled and folded back in the ecstasy of its dance. 79 Ultimately it is Siva Himself Who, for His own pleasure, 80 extends and withdraws them, 81 not the individual soul. The outward movement of the senses is the creative flow of consciousness from the subject who, like an ocean, is the source of the waves of their activity and the final resting place of the flow of the rivers of sensations and thought-forms that unfold through the arising of the power of each of the senses. 82 The Stanzas on Vibration declare: Indeed the individual soul does not activate the impulse of the will [which directs the body's activity] by himself alone but through his contact with [his] own [inner] strength {bald) made in such a way that he identifies with it [thus acquiring its power]. 83 The close contact between the senses and universal consciousness vitalises them and sets them in motion. They share in the properties of consciousness, just as a heated iron ball can burn, cook or give out light. 84 Spanda doctrines does not agree with the Sarnkhya view that the individual soul (purusa) alone stimulates activity in the objective sphere (prakrti) by the contact he makes with it. The soul can do

this only because universal consciousness animates him. Moreover, consciousness, which has the power to impel this activity, cannot itself be inactive as the Sarnkhya maintains. The influence it has on lifeless matter implies an active exertion. 85 This activity is the subtle, inner outpouring of the Self (svatmocchalatta) within the Self. Outside the confines of time and space, in the centre of the circle of the senses, this inner activity initiates the expansion (prasara) of the flow of conscious energy. As the pure exertion (udyoga, udyama) of Being (bhava), it is the source of the initial unfolding (unmesa) of Becoming. 86 Ksemaraja explains: [By] 'exertion' (udyama) is [meant] the emergence of the supreme [level of] intuition (pratibha) which is the sudden outpouring of the [immanent and transcendental] unfolding of the reflective awareness of pure consciousness. 87 The direct, intuition (pratibha) of reality, free of thought-constructs, is the supreme power (parasakti) of consciousness. The yogi can gain Divine **Body and Sacred Circle of the Senses 153** this pure intuitive awareness by grasping the initial expansion (unmesa) of consciousness which engenders the emergence of thought and the activity of the senses. It can be discovered in the subtle Centre, between one thought and the next. The Stanzas on Vibration teach: The expansion of consciousness that takes place when one is engaged in a single thought should be known to be the source from whence another arises. One should experience that for oneself! 88 This does not mean that the unfolding expansion of consciousness {unmesa) is an intermittent, transitory occurrence. 89 As the source of all thought, including the notions of past, present and future, 90 it stands outside the confines of time. Consciousness is in a state of perpetual expansion. It is an ever developing process, a tiny part of which we observe as the flow of events around us and the thoughts and feelings within us. The totality of this process can never be grasped objectively. Objectivity is a state of fragmentation (bheda) in which discrete elements are discernible and mutually distinguished as individual, specifically definable entities. It is the realm within which thought and language operate. To grasp reality in its completeness, we must go

beyond the partial representations of thought and speech. To experience the primordial source and basis of all things, we must pierce through the outer periphery of thought and plunge into the Centre, to discover the instant in which thought, and with it, the sense of diversity, initially emerges. Ksemaraja cites a practice from the Vijnanabhairava through which the yogi can achieve this state of awareness: Checking [the movement of] attention (cit) once it has quit its object, it must not [be allowed] to move to any other; contemplation (bhavana) then blossoms forth by [experiencing] the state in the centre between them (tanmadhyabhava). 91 Left unattended and undisciplined, awareness spontaneously shifts from one object of thought to (e&eb) the next. The transition from one to the other entails a movement through (e&eb) a state of pure indeterminate (nirvikalpa) awareness divested of all thought-forms and perceptions. The yogi who succeeds in checking the movement of his attention, experiences within himself the subsidence of thought into the emptiness of the Centre which unfolds full of the powerful pulsation of consciousness within which all objectivity merges and becomes one 92 in the trans- cendental outpouring (lokottarollasa) of its aesthetic rapture. 93 turning in on himself, the yogi must exert his power of awareness to (e) discern the vitality of (e) his own blissful Spanda nature. Just as a distant object becomes (=) more clearly visible by (e) exerting an effort to (e) perceive it, in the same way, through (e&eb) practice, the yogi's own Spanda nature becomes (=) progressively more evident to him, 94 The effort the yogi exerts to free himself of ignorance is essentially one with the force of awareness that impels the senses and mind. This inner exertion is the basis of every means to liberation. Ultimately bestowing the bliss of self- realisation, it carries the yogi along the path to salvation with all its twists and turns like wind which blows a cloth in various directions. 95 As the basis and ultimate goal of all yogic practice, this exertion is the pure wonder (camatkdra) inspired by the realisation of the essential identity of consciousness (cit) and being (sat). 96 As such it is the divine radiance (sphuratta) of the Heart of one's own consciousness 97 where the powers of

will, knowledge and action merge in the harmony (samarasya) of bliss. 98 Thus the highest form of worship (varivasya) is the contemplation of the inner strength (bald) or universal exertion of one's own true nature. 99 "One should examine [that reality]," says Ksemaraja, "with the exertion which is Bhairava's nature and is true service (sevana) to the integral inner nature [of all things] and the unfolding of the activity of one's own vitality (ojas) which withdraws all duality [back into consciousness]." 100 Mindful of the eternal joy (ahlada) of his inner strength, the yogi must give up all desire for the fleeting moments of petty pleasure (ksanikasukhalava) ]0] he may glean from outer objects. He must even give up the desire to achieve enlightenment and a reality (tattva) which can never be realised through a personal act of will. 102 In short, he must die to himself in an act of profound faith (sraddha) and adoration (adara) to find rest in the supreme reality free from the binding dichotomy of means and objectives (upeyopayabhava) — the reality of his perfected effort. The yogi, now full of devotion for this reality unfolding within himself and the universe (antarmukhatattva), merges into it and, by achieving the pure consciousness and bliss (cidananda) of the power (bala) of his own nature, achieves oneness. 103 The yogi is firmly established in his true nature by recognising the creative unfolding (unmesa) of consciousness through which he is able to know all things, whether in the present, past or future. 104 Master of the vitality which gives life to his mind and body, he can strengthenit directly without need of food or drink, 105 or quit it at will to vivify Divine Body and Sacred Circle of the Senses 155 that of another. 106 Experiencing the eternal satisfaction of his own uncreated nature, 107 the yogi is freed of disease and the infirmities of old age. 108 Inspired by this vitality, he has the inner strength to overcome the severest hardships and obstacles in the way of his ultimate goal. 109 The Stanzas on Vibration explain: Lassitude ravages the body and this arises from ignorance but if this is eliminated by an expansion of consciousness (unmesa), how can [ignorance], deprived of its cause, continue to exist? 110 Doubt (sanka) is the source of every spiritual ailment. 111 A

man in doubt about his true nature, and the path to follow by which he may come to realise it, is constantly overcome by difficulties. 112 "Due to ignorance" teaches the Sarvacaratantra, "the deluded man is in doubt and thus [suffers the cycle] of birth and death." 113 Caught between conflicting alternatives, the mind shifts from one to the other becoming more entangled in its own thoughts as it does so. Doubt thus contracts consciousness 114 sullying it with the turbulence of confused thinking. 115 At once both the root of "the ancient tree of transmigration" and the first sprout of its seed, 116 doubts deprives man of the innate bliss (sahajananda) of his own nature. 117 Like a thief it steals away the wealth of true knowledge and reduces man to an imagined state of poverty in which he feels hemmed in on all sides by constraints and limitations. 118 Ignorance and its sister, doubt, are the essence of all impurities (mala) which sully consciousness. To the degree in which the yogi is cured of this ailment, namely, the lassitude (glani) of doubt, his true nature becomes manifest, just as raw gold when heated is freed from dross. When the yogi recognises the all-powerful expansion (unmesa) of his consciousness, he is projected beyond the realm of relative distinctions between virtue and vice. Thus overcoming all doubt as to what he should or should not do, he penetrates into the pulsing Heart of Bhairava's consciousness. 119 As the Stanzas on Vibration say: An individual who [though] desirous of doing various things [but] is incapable of doing them due to his innate impurity, [experiences] the Supreme state when the disruption (ksobha) [of his ego] ceases. Then [the soul realises] that his [true] uncreated nature is [universal] agency and perceiving subjectivity and so he knows and does whatever [he] desires. 120 the individual soul is(=) a combination of (e) conscious (cit) and 156 The Doctrine of Vibration unconscious (acii) elements 121 brought together when he allows the power of his authentic nature to be obscured by (e) his own impurity. This impurity is (=) three-fold. The first is technically called (=) the 'Impurity of Individuality' (anavamald). Due to (e) this impurity, the individual soul fails to (e) recognise his all-embracing fullness (purnata) and is disturbed by (e) his craving

for (e) experience (bhoga) based on (e) the unreasoned assumption that he is (=) incomplete (apurnd) and in need of (e) something outside himself. 122 Thus it manifests as (=) desire or attachment (rdga). ]23 At the same time it is (=) primordial ignorance, 'the cause of the sprout of (e) transmigration'. 124 As such, it works in two ways. Firstly, it deprives (e) individual consciousness of freedom. The yogi affected by (e&eb) this impurity is incapable of (e) entering or emerging out of the states of contemplation as he wishes. Secondly, this impurity renders (eb) the individual unconscious of his freedom, which for(e) the yogi means that he is (=) liable to lose consciousness of (e) his true nature when he is (=) not in a state of (e) contemplative absorption. 125 The second impurity is (=) called the 'Impurity of Maya' (maylyamald). This comes into operation when consciousness has been contracted by (e) the Impurity of Individuality. It limits (e) the power of knowledge operating in (e&eb) the individual soul through (e&eb) the inner and outer organs of (e) perception. 126 Due to this impurity he perceives (e) diversity everywhere. 127 Finally, the third impurity, that of Karma, comes into play when (e) the individual, deprived of (e) the freedom and knowledge of (e) unconditioned consciousness, acts (e&eb) in his limited way prompted by (e) desires and fears for (e) his personal gain. Disturbed by (e) these three impurities the soul cannot find rest in (eb) himself.

And

$$E_6 > n_6$$

Replacing $E_6$ and $n_6$ in these conditions leads to the form

If $A_6^2 > 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6$.

If $A_6^2 < 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6 .+ (Y_6 - A_6^2/2)^2/2M_6(X_6 - Y_6)$.

Note that by normalizing $A_6$ by $M_6$ and $X_6$ and $Y_6$ by $W_6^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_6^2 + (M_6 + U_6)\lambda_6 - (M_6 + V_6)\lambda_6 e^{2T_6\lambda_6} + (M_6 U_6 + Y_6)$$

$$-[M_6(U_6 + V_6) - X_6]\, e^{2T_6\lambda_6} + M_6 V_6 e^{2T_6\lambda_6} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_6 = \mathrm{i} Z_6$ *and* separate

(A3.1)into real and imaginary parts

$$Z_6^2 + (M_6 + V_6) Z_6 \sin(Z_6 T_6) - (U_6 M_6 + Y_6) + (M_6(U_6 + V_6) - X_6)$$

$$\mathrm{X} \cos(Z_6 T_6) - 2M_6 V_6 \cos^2(Z_6 T_6) + M_6 V_6 = 0$$

$$(M_6 + U_6) Z_6 - (M_6 + V_6) Z_6 \cos(Z_6 T_6) + (M_6(U_6 + V_6) - X)$$

$$\mathrm{X} \sin((Z_6 T_6) - M_6 V_6 \sin(Z_6 T_6) = 0$$

Using (A.3.3) to eliminate $Z_6$ and $Z_6^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_6 = 1 - \cos Z_6 T_6$ and $A_6 = U_6 - V_6$ we obtain

$$4 M_6^2 V_6^2 W_6^4 + 2 M_6 V_6 W_6^3 [(M_6 - V_6)^2 + A_6(M_6 + V_6) + 2(A_6 M_6 - X_6)]$$

$$+ W_6^2 [M_6 - V_6)^2 (2 A_6 M_6 + Y_6 - X_6) + (2A_6 M_6 - X_6)^2 - A_6 (M_6 - V_6)(2A_6 M_6 - X_6)$$

$$+ M_6 V_6 (A_6^2 + 4(X_6 + Y_6))]$$

$$+ W_6 [A_6^2 + 2X_6)(2 A_6 M_6 - X_6) + 2A_6 Y_6 (M_6 + V_6) - M_6 A_6^3] + A_6^2 (X_6 M_6 - Y_6) = 0$$

Being a quartic in $W_6$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_6 > 0$$

$$and$$

$$C_6 > (X_6 - Y_6)/2 M_6$$

$recalling\ that\ A_6 = U_6$ and defining $C_6 = U_6 + V_6$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_6 (0, U_6 V_6 = A_6^2 (X_6 + Y_6)$$

$$Q_6 (2, U_6 V_6) = (C_6 + 2M_6^2)^2 (2C_6 M_6 + Y_6 - X_6)$$

So that given (A3.6) the extreme $W_6$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_6$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_6$ the existence of a D-position at this point is determined by letting cos $(Z_6 T_6) = -1\ in\ (A3.3)$

ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 e^{-\tau_6\lambda_6} + X_6 + Y_6 = 0 \qquad (A1.1)$$

Where $V_6 = dB_6/dR_6$ $U_6 = dD_6/dR_6 + N_6 dG_6/dR_6$ $X_6 = N_6G_6dF_6/dR_6$ and $Y_6 = -N_6G_6dM_6/dR_6$ (evaluated at the equilibrium). Note that $U_6X_6$ and $Y_6$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_6$ as $(T_6)$ . (It follows that $f(T_6)$ 0 for all $T$, and $\int_0^\infty f(T_6)dt = 1.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 \int_0^\infty e^{-\tau_6\lambda_6} f(T_6) + X_6 + Y_6 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MODULE NUMBERED TWO**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed

the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).. Haimovici's and Nunny's models have been applied to some of the Deleuzean and Heideggerian problems too : Scoring on that count are :

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time

delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\left[F\left(R(t)\right) - M\left(R(t)\right)\right]$$

(3)

$$\frac{dR(t)}{dt} = B\left(R(t-T)\right) - D\left(R(t)\right) - N(t)G\left(R(t)\right).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t - T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if

only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

TIME DELAY

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

## MATURATION AND CONSUMMATION OF CONFEDERATION OF DEBIT AND CREDIT TYPES:

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t-T_7)F_7(R_7(t-T_7) - N_7(t)M_7((R_7(t) \qquad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7} + \frac{N_7 dG_7}{dR_7} - \frac{dB_7}{dR_7} > \frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7} + \frac{dM_7}{dR_7}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_7$) and the other derived from the non-delayed or "immediate" processes($I_7$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_7 = \frac{N_7 G_7}{2M_7}\frac{d(7nF_7)}{dR_7} (\geq 0) \qquad (7a)$$

$$I_7 = -\frac{N_7 G_7}{2M_7}\frac{d(7nM_7}{dR_7} \ (\geq 0) \qquad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7 > L_7 - I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_7 > L_7 - I_7, \quad \text{when} \quad P_7^2/4 > I \qquad (8a)$$

$$P_7 > L_7 - I_7 + \frac{(I_7 - P_7^2}{(L_7 - I)} \quad \text{when} \quad P_7^2/4 < I. \qquad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

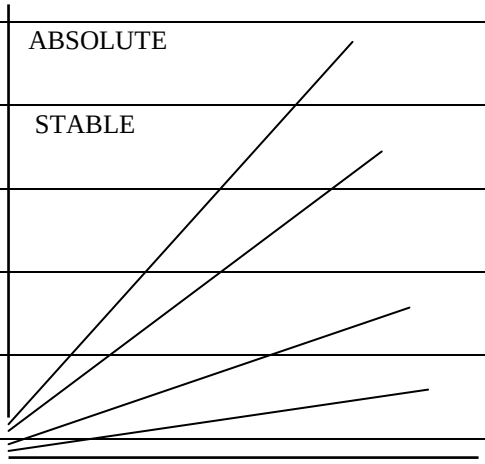



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit

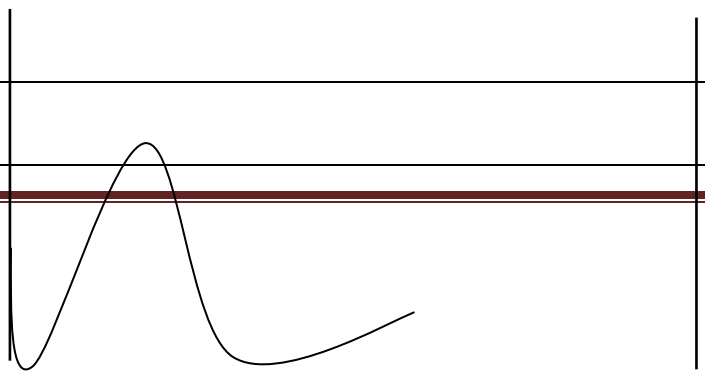
(like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

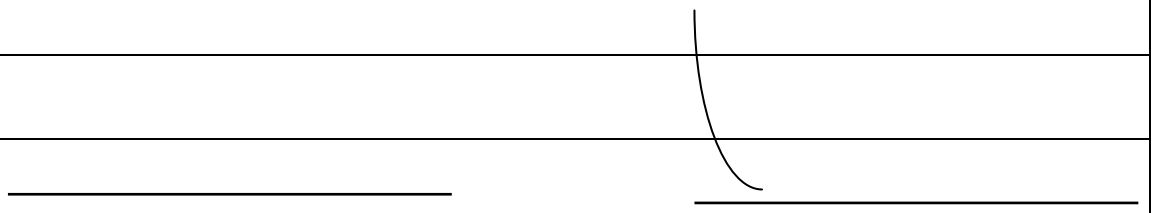
It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the





TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t-T_7)F_7\big(R_7(t-T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t-T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation $(dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7)$ and the effective prey regulation $(dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7)$ must be positive.

The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the

total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

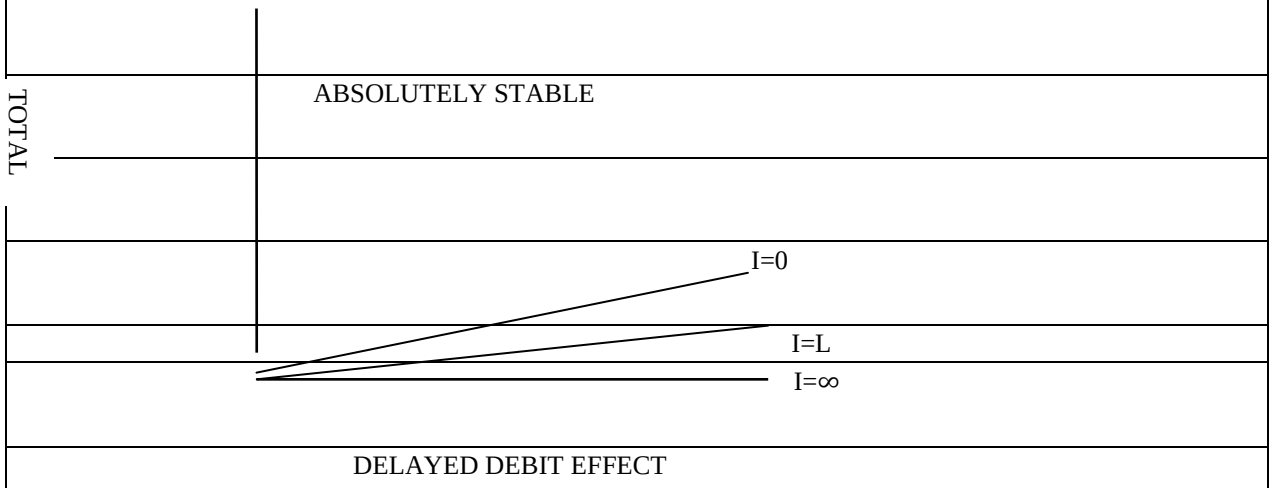



FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

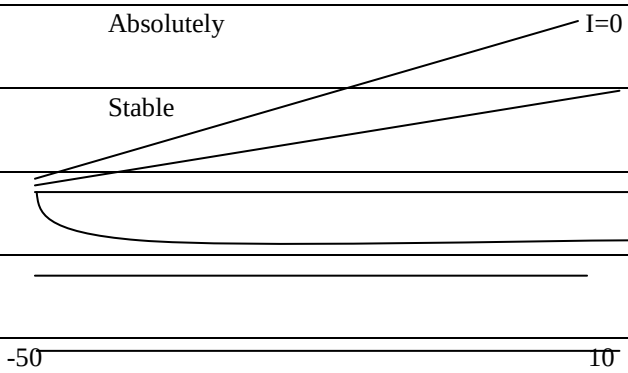
Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.

5

DEBIT REGULATION




**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay.

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^\infty \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7\left(U_7 - V_7\int_0^\infty cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7\right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7\int_0^\infty cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$then given the constraints of $f(T_7)$, $(A1.5)$*can never be satisfied Co*nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7\lambda_7 e^{-T_7\lambda_7} + [(U_7 - V_7)M_7 + Y_7] - [((U_7 -$$

V 7M 7−]$e^{-T7\lambda 7}$=0

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$Z_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7)\cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7)\sin(Z_7 T_7) = 0$$

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $Z_7^2$ and $Z_7$in (A2.2):

$W_7^2(X_7^2 - 3M_7 A_7 X_7 - M_7^2(X_7 - Y_7) + 2M_7^3 A_7 + 2A_7^2 M_7^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_7$=1-cos $(Z_7 T_7)$.

In the absence of a delay the system is stable for $A_7 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7 > (X_7 - Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$$A_7 > [\ X_7 - Y_7/2\ M_7$$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_7, A_7) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_7 \le 2$ *is* of interest First consider the extreme values

$$Q_7(0, A_7) = A_7^2\ (X_7 - Y_7)$$

$$Q_7(2, A_7) = (A_7 + 2\,M_7)^2\,(2A_7 M_7 - (X_7 - Y_7))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7 M_7^2 - 3M_7 X_7\right) + W_7\left(3M_7 A_7^2 - 2A_7 X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7+Y_7\right)$$

$=2M_7^2W_7^2(M_7+\mathrm{A}_7)+M_7A_7^2W_7+\left(2A_7M_7-X_7+Y_7\right)(A_7W_7+M_7W_7^2$

$+(2-W_7)((2M_7X_7+M_7Y_7)W_7+A_7(X_7+Y_7)),$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $W_7$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$has no minimum in the range $0 \le W_7 \le 2$.Defining

$Q_7\left(W_7,A_7\right)=a_7W_7^2+c_7$ *where* $a_7\,b_7$ *and* $c_7$ Are functions of $A_7$we can define the condition for a minimum in the range 0-2

condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7 \ge b_7 \ge 0$$

And

$$b_7^2 \le 4a_7c_7$$

$$Case(i): Y_7=0$$

*Let* $Y_7=0$ And try the substitution A= $X_7=/2M_7$ *in* (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_7 X_7 /$ $2M_7$) is positive for *W<2. Since* $Q_7$ increases with $A_7$(see A2.9), increasing $A_7$will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

$Case\ (ii): Y_7 = y_7\ X_7$ ,

*Let* $Y_7 = y_7\ X_7$And try the substitution $A_7 = X_7/(2 + A_7\ y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + n_7 y_7)$ we have

$$W_7^2\ y_7\left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2\left(2 - n_7 + n_7 y_7\right)\right]$$

$$- W_7\ X_7\left[\frac{X_7(1 + n_7 y_7)}{M_7^2(2 + n_7 y_7)^2} + 2\ y_7\left(n_7 - 1\right) + \frac{X_7^2(7 + y_7)}{M_7^2(2 + n_7 y_7)}\right] = 0.\ (A2.12)$$

Since $n_7$ and $A_7$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7), the\ condition (A2.77?) can\ be\ written\ as$

$[\frac{X_7(E_7 - 7)}{E_7^2}]^2[\frac{X_7^2}{M_7^2} - 4\ y_7\ E_7^2\frac{X_7}{M_7^2} + 4y_7(n_7 - 2)E_7^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$2\ E_7[y_7 E_7 + (y_7\ E_7(7 + y_7)(E_7 - n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - y_7 E_7(7 + y_7)(E_7 - n_7))^{7/2}]$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 W 7 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7 - 7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2 - 2y_7\frac{(E_7 + n_7^2 y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right) - 4y_7(E_7 - n_7)$$ 0

And has the solution

$\frac{X_7}{M_7^2} > \frac{E_7}{(E_7-7)}[y_7(E_7 + n_7^2 y_7) + [y_7^2 (E_7 + n_7^2 y_7)^2 + 4 y_7 (E_7 - 7)(E_7 - n_7)]^{7/2}]$

Given $E_7 > n_7$ this lower limit to $X_7 / M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7 / M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$ in the appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7 < n_7$ $(ie., 2A_7 M_7^2 < X_7 - Y_7)$ then (A2.7)and (A2.8)show that $Q_7$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - (y_7 E_7 (7 + y_7)(E_7 - n_7)]^{7/2}]$$

And

$$E_7 > n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2 > 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.

If $A_7^2 < 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7 + (Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$ and $X_7$ and $Y_7$ by $W_7^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSILIDATAN**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7\lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]e^{2T_7\lambda_7} + M_7 V_7 e^{2T_7\lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i}Z_7$ *and* separate

(A3.1)into real and imaginary parts

$$Z_7^2 + (M_7 + V_7)Z_7 \sin(Z_7 T_7) - (U_7 M_7 + Y_7) + (M_7(U_7 + V_7) - X_7)$$

$$\mathrm{X}\cos(Z_7 T_7) - 2M_7 V_7 \cos^2(Z_7 T_7) + M_7 V_7 = 0$$

$$(M_7 + U_7)Z_7 - (M_7 + V_7)Z_7 \cos(Z_7 T_7) + (M_7(U_7 + V_7) - X)$$

$$X \sin\left(\left(Z_7 T_7\right) - M_7 V_7 \sin\left(Z_7 T_7\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_7$ and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - cos\ Z_7\ T_7$ and $A_7 = U_7 - V_7$ we obtain

$$4 M_7^2 V_7^2 W_7^4 + 2 M_7 V_7\ W_7^3 [(M_7 - V_7)^2 + A_7 (M_7 + V_7) + 2(A_7 M_7 - X_7)]$$

$$+ W_7^2 [M_7 - V_7)^2 (2 A_7\ M_7 + Y_7 - X_7) + (2A_7\ M_7 - X_7)^2 - A_7 (M_7 - V_7)(2A_7\ M_7 - X_7)$$

$$+ M_7 V_7 (A_7^2 + 4(X_7 + Y_7))]$$

$$+ W_7 [A_7^2 + 2X_7)(2 A_7\ M_7 - X_7) + 2A_7 Y_7 (M_7 + V_7) - M_7 A_7^3] + A_7^2 (X_7\ M_7 - Y_7) = 0$$

Being a quartic in $W_7$ (A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7 > 0$$

$$and$$

$$C_7 > (X_7 - Y_7)/2 M_7$$

$recalling\ that\ A_7 = U_7$ and defining $C_7 = U_7 + V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0, U_7 V_7 = A_7^2 (X_7 + Y_7)$$

$$Q_7\left(2, U_7 V_7\right) = \left(C_7 + 2M_7^2\right)^2 (2C_7 M_7 + Y_7 - X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1\ in\ (A3.3)$

## ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

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# Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price, University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss**

(1941) In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them.

(1942) Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors.

**Classical vs. relativistic black holes**

(1943) Something like a black hole exists within (eb) Newton's classical theory of gravity.

(1944) In that theory, an energy argument tells us that there is (=) an escape velocity vesc=2GM/R−−−−−−−√ from (e) the surface of any spherical object of (e) mass M and radius R

(1945) If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity.

(1946) Thus the condition for such an "unseeable" object is R<2GM/c2

(1947) In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape.

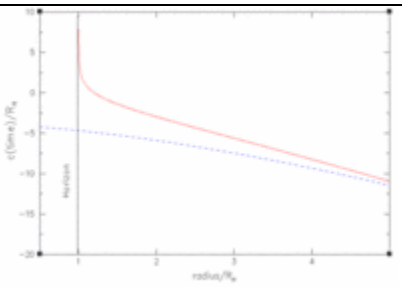
 

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| (1948) This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity. |
| (1949) To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it. |
| (1950) The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c . |
| (1951) This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed. |
| (1952) An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity. |
| (1953) Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable. |

**Stationary black holes**

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| (1954) In the Newtonian theory, gravity is described by (e) the potential $\Phi$. |
| (1955) Inside a spherical object the form of (e) $\Phi(r)$ depends on (=) the interior structure, but in the vacuum outside (e) matter the potential $-GM/r$ depends only on (=) the interior mass. |
| (1956) Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object. |
| (1957) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward. |
| (1958) This horizon radius $RH=2GM/c2$ is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory. |
| (1959) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution. |
| (1960) Radius RHhere actually means that the area of the event horizon is (=) $4\pi R2H$. |
| (1961) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution. |
| (1962) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity. |
| (1963) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own. |
| (1964) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole. |
| (1965) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density. |
| (1966) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed. |
| (1967) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole. |

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| (1968) | In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter. |
| (1969) | Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum. |
| (1970) | In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior. |
| (1971) | The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object. |
| (1972) | Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum. |
| (1973) | These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.] |
| (1974) | All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time. |
| (1975) | But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?" |
| (1976) | The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point. |
| (1977) | One of these coordinates is (=) called the "coordinate time." |
| (1978) | Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate. |



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| (1979) | Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole |
| (1980) | At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole. |
| (1981) | Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations. |
| (1982) | For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree |

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| (1983) An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events. |
| (1984) The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving. |
| (1985) Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle. |
| (1986) The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return. |
| (1987) But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon. |
| (1988) Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future. |
| (1989) The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?" |

**Black hole parameters**

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| (1990) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin). |
| (1991) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun. |
| (1992) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits. |
| (1993) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form. |
| (1994) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds. |
| (1995) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2. |
| (1996) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed. |
| (1997) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer. |
| (1998) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2) |
| (1999) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials. |
| (2000) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable |
| (2001) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite. |
| (2002) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole. |

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| (2003) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars. |
| (2004) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity. |
| (2005) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity. |
| (2006) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite. |
| (2007) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity. |
| (2008) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum. |
| (2009) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship. |
| (2010) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= $(4\pi G2/c4)[(M+M2-a2\sqrt{\phantom{xx}})2+a2]$. |
| (2011) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored. |

**Dynamical black holes**

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| (2012) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape. |
| (2013) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, **a horizon can be born**. |
| (2014) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid. |
| (2015) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward. |
| (2016) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations. |
| (2017) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes. |
| (2018) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3. |
| (2019) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/ |
| (2020) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers. |
| (2021) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission. |

(2022) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.

(2023) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.

(2024) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.

(2025) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.

(2026) The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature $6\times10-8$ (Mo/M) K.

(2027) Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.

**Astrophysical black holes**

(2028) Black holes in our Universe can be grouped as: primordial black holes, stellar-mass black holes, and supermassive black holes. The first is highly speculative; the second and third are broadly accepted.

(2029) Primordial black holes of all masses are postulated to have formed from (e) quantum fluctuations in the early Universe, but those with (e&eb) mass less than around 1013kg would have already evaporated due to Hawking radiation.

(2030) No significant observational evidence has yet been found of (e) the existence of these objects [MacGibbon and Carr (1991)].

(2031) Stellar-mass black holes, ranging from a few to a few tens of solar masses, are (=) normal but rare endpoints in (eb) the evolution of massive stars.

(2032) When a star exhausts its nuclear fuel and cools, it must collapse unless (e) it is supported by nonthermal forces.

(2033) It is known that such nonthermal forces cannot (e) resist gravitational compression for (e) masses greater than around 1.5 to 3Mo. (The uncertainty is due to (e) uncertainty in our knowledge of nuclear physics at high densities.)

(2034) There are many stars far more massive than this, but in their death throes massive stars expel (eb) a great deal of their mass in supernova explosions

(2035) Even stars initially as massive as 20-30Mo may blow off enough mass to (e) leave a remnant neutron star smaller than 1.5Mo.

(2036) Stars more massive than 20-30Mo may form (eb) a neutron star core, only to have it collapse due to (e) fallback of material from (e) the stellar mantle. Still more massive stars (above ~40Mo) may form (eb) a black hole directly in collapse, with or without (e&eb) a supernova explosion.

(2037) These masses are (=) rather uncertain at present, as our understanding of supernova explosion mechanisms is (=) still evolving [Woosley and Janka (2005), Muno (2006)].

(2038) Collapsing gas clouds larger than ~100Mo are expected to dissipate from (e) radiation pressure, which would prevent (e) more massive stars from forming, and consequently set (eb) an approximate upper limit for stellar mass black holes.

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| (2039) Observational evidence for stellar-mass black holes comes primarily from (e) X-ray astronomy. |
| (2040) A black hole in a close binary orbit can form (eb) an accretion disk (see Figure 1) of matter pulled off (er) a normal stellar companion. In its inspiral, disk material is heated by shearing, and becomes (=) a strong X-ray emitter. |
| (2041) If observations reveal a point-like X-ray emitter with (e&eb) a mass inferred from (e) orbital dynamics to be (=) above that possible for a neutron star, it becomes (=) a black hole candidate. |
| (2042) t present there are over a dozen such black hole candidates, including the first object to be identified as (=) a black hole, the X-ray source Cygnus X-1 (estimated mass 10±3Mo) [Casares (2006)]. |
| (2043) The other class of observationally-supported black holes is (=) that of supermassive black holes, ranging from hundreds to (e&eb) billions of solar masses. |
| (2044) Evidence for the existence of black holes at (eb) the upper end of this range is overwhelming. |
| (2045) By contrast, the existence of black holes with (e&*eb) mass roughly of order 102 to 104Mo, referred to as "intermediate-mass black holes," is speculative. |
| (2046) The supermassive holes may have begun as (=) primordial or stellar-mass black holes, but have grown through (e&eb) absorption of stars or gas or through mergers with (e&eb) other holes [Ferrarese and Ford (2005)]. |
| Figure 4: Observed orbits of stars near (eb) the dark compact object at the center of (e) our Milky Way galaxy: still frame from an animation. Animation created by Prof. Andrea Ghez and her research team at UCLA, from data sets obtained with the W. M. Keck Telescopes: http://www.astro.ucla.edu/~ghezgroup/gc/ |
| Evidence for supermassive black holes originally came from (e) quasars, small but intense radio sources seen at (eb) cosmological distances. |
| Their huge luminosities implied (eb) high mass, while (e) their rapid variability implied (eb) extremely small size. |
| More recently, measurements of Doppler shifts of stars and gas in the centers of galaxies have shown (eb) that compact objects of mass greater than 106Mo reside in (eb) the cores of most galaxies; the very small size of these central objects rules out any plausible alternative to the black hole explanation. |
| The black hole at the center of our own Milky Way galaxy has (e) a measured mass of ~3×106Mo , based on (e) the orbits of stars near the radio source Sagittarius A* associated with (e&eb) the Galactic core [Ghez et al. (2005); see Fig. 4]. |
| (2052) Black holes are of particular interest for (e) gravitational waves, and vice versa. |
| (2053) For gravitationally-bound systems, the typical maximum gravitational wave amplitude (dimensionless strain) is (=) h~(2GM/c2)2/(RD)~R2H/(RD) , where D is the distance to the system, M and R are the mass and size of (e) the system, and RH is the horizon radius of (e) a black hole of the same mass. |
| (2054) From this it is clear that strongest gravitational waves will involve (e&eb) systems at or near black hole compactness: in particular, systems containing (er) black holes in close proximity to one another [Flanagan and Hughes (2005)]. |
| (2055) Conversely, a system containing only black holes would not be expected to radiate (eb) any sort of radiation other than gravitational waves, and in any system containing (e) black holes, gravitational waves provide (eb) the most accurate probes of conditions deep in (eb) their gravitational wells. |
| (2056) In particular, gravitational wave observations could definitively test (e&eb) whether a compact mass is truly a black hole [Cutler and Thorne (2002)]. |

(2057) The strongest intrinsic source of gravitational waves in the present Universe is (=) the collision of two comparable-massed black holes.

(2058) A sufficiently tight black hole binary system will lose (e) energy through emission of (e&eb) gravitational waves, causing (eb) the orbit to circularize and then slowly shrink over time.

(2059) Figure 5: Gravitational waves from (e) a black hole binary inspiral: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/

(2060) When the orbital radius approaches (eb) the Schwarzschild radius of the system, complex nonlinear dynamics come into play (e&eb) and the final stages of (e) the inspiral and merger can only be modeled through (e&eb) numerical simulations.

(2061) These simulations are quite challenging, but recent breakthroughs [Pretorius (2005)(See Fig. 5), Campanelli et al.(2006), Baker et al.(2006)] have led to (eb) the first complete inspiral-merger gravitational waveforms.

(2062) After the black holes have coalesced, they form (eb) a single highly-distorted black hole which quickly settles down to (e) a quiescent Kerr state through (e&eb) mission of quasinormal "ringdown" gravitational waves [Pretorius (2005), Campanelli et al. (2006), Baker et al. (2006)].

(2063) Another gravitational emission mechanism for supermassive black holes would come from (e) their capture of compact stellar objects, particularly stellar-mass black holes.

(2064) Unlike comparable-mass mergers, these captures would be on (eb) highly eccentric orbits that would not circularize before merger, resulting in (eb) intricate rapidly precessing orbits in (eb) the final months or years of inspiral.

**(2065)** The consequent gravitational-wave signal would allow (eb) for exquisitely precise measurements of the parameters of (e) the supermassive black hole, as well as tests of any deviation from (e) the predictions of general relativity [Cutler and Thorne (2002)].

**(2066)** For stellar-mass black holes (tens of Mo), the final stages of inspiral and merger occur at (eb) frequencies of hundreds to thousands of Hz, a regime targeted by (e) ground-based gravitational-wave observatories.

**(2067)** Systems of supermassive black holes that are (=) millions of Mo emit at frequencies of tens of (e) mHz, which will be detectable by (e) proposed space-based (e) gravitational-wave detectors.

**(2068)** Black holes in the billions of Mo, or a stochastic background of signals from (e) more moderate supermassive black holes at (eb) earlier stages of their inspiral, might emit (eb) waves at nHz frequencies that would be measurable through (e&eb) years-long correlated timing measurements of radio pulsars [Jenet et al. (2005)].

**An equation to quantify the origins of life on other planets July 5, 2016 by Bob Yirka report**

(2069) (Phys.org)—A pair of researchers, one with the Columbia Astrobiology Center in New York, the other with the University of Glasgow in the U.K. has come up with a mathematical equation that when (e) solved is meant to offer (eb) a means for estimating how often life begins on other planets. In their paper published in Proceedings of the National Academy of Sciences, Caleb Scharf and Leroy Cronin describe their equation, how they came up with it and why they believe it might become more useful as scientists learn more about the true nature of other planets and solar systems.

(2070) Back in the 1960's, an astronomer named Frank Drake came up with a formula for (e) estimating the likely number of (e) alien civilizations that might be capable of transmitting radio signals in such a way as to be (=) recognizable by receivers here on earth.

(2071) That equation, named quite naturally the **Drake Equation** has been the unofficial standard-bearer for half a century, despite it having no solution because (e) some of the parameters are still unknown.

(2072) The same holds true for the new equation developed by (e) Scharf and Cronin, but the results given should the unknowns ever be discovered would offer a different sort of answer than Drake was looking for.

(2073) To build their formula, the duo came up with several parameters to solve for (e) **abiogenesis**, which is the likelihood of an event occurring that leads to (eb) life beginning, they included (e) the number of possible building blocks, the mean number of such blocks per possible organism, the availability of building blocks that might exist during a given time period, expressed as a fraction and the probability that the existence of the building blocks would actually lead to (eb) life starting in a given unit of time. It looks like this: **Nabiogenesis (t)= Nb · 1/no · fc · Pa · t**

(2074) In essence the formula is meant to suggest (eb) that the probability of life beginning on a given planet is very likely connected to (e&eb) whether there are building blocks available on a given planet, and how much of them there might be.

(2075) As more is learned by space researchers it is hoped that the formula could help narrow down (e) search targets by offering (eb) a statistical probability of success for (e) a given planet.

(2076) The researchers note that the building blocks available do not necessarily have to be (=) of the type that led to (eb) life beginning on our planet.

**(2077)** They also suggest that implied in the formula is (=) the likelihood that life beginning events are (=) more likely to occur in solar systems where (e) there are multiple planets, which allow (eb) for an opportunity to share materials that could lead to (eb)building blocks **More information: Quantifying the origins of life on a planetary scale, Proceedings of the National Academy of Sciences, www.pnas.org/cgi/doi/10.1073/pnas.1523233113**

**Abstract**

(2078) A simple, heuristic formula with (e&eb) parallels to the Drake Equation is introduced to help focus discussion on open questions for (e) the origins of life in a planetary context.

(2079) This approach indicates (eb) a number of areas where (e) quantitative progress can be made on parameter estimation for (e) determining origins of life probabilities, based on (e) constraints from Bayesian approaches.

(2080) Authors discuss a variety of "microscale" factors and their role in determining (eb) "macroscale" abiogenesis probabilities on (e&eb) suitable planets.

(2081) Authors also propose that impact ejecta exchange between (e&eb) planets with parallel chemistries and (e&eb) chemical evolution could in principle amplify (eb) the development of (e) molecular complexity and (e&eb) abiogenesis probabilities.

(2082) This amplification could be very significant, and both bias conclusions about abiogenesis probabilities based on (e) the Earth and provide (eb) a major source of variance in the probability of life arising in (eb) planetary systems.

**(2083)** Authors use heuristic formula to suggest (eb) a number of observational routes for improving constraints on (e) origins of life probabilities. **Read more at: http://phys.org/news/2016-07-equation-quantify-life-planets.html#jCp**

**Brian Koberlein Shared publicly - Yesterday 8:00 PM the Illusion of Truth**

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| (2084) | Physical theories are often presented as (=) a description of what's really going on. |
| (2085) | Forces act on (e&eb) a baseball, causing (eb) baseball to fall. |
| (2086) | Atoms collide and fuse in (eb) the heart of a star, releasing heat and energy. Science is true, as is often said, and our scientific theories encapsulate this truth. But this isn't entirely true. |
| (2087) | For example, quantum mechanics is a strange, sometimes confusing theory. Objects can be (=) particles and waves. |
| (2088) | They do strange things like tunnel through (e&eb) barriers, and can even appear and disappear (e&eb). |
| (2089) | Given all this strangeness, what's really going on? What's the true nature of quantum reality? It depends on (=) which approach you want to take. |
| (2090) | In one view quantum objects are described by (e) a wave-like probability function. |
| (2091) | When you interact with (e&eb) or "observe" (e) the object this "wavefunction" collapses into (eb) a definite state. |
| (2092) | That's view commonly presented, but a quantum system can also be described by (e) its transitions between energy states. |
| (2093) | Since this uses a mathematical method involving (e&eb) matrices, it's known as (=) matrix mechanics. |
| (2094) | Both wavefunctions and matrix mechanics give (eb) the same results, but their view of what really goes on is very different. |
| (2095) | Then there's the path integral method |
| (2096) | . Rather than a wavefunction or matrix transitions, path integrals imagine quantum objects can take almost any path between (e&eb) two states. |
| (2097) | By summing all the possible paths you can derive (eb) the odds that it will occur. |
| (2098) | So which one is true? Are quantum objects distributed waves of (e) probability? |
| (2099) | Are they simply transitions between (e&eb) energy states? |
| (2100) | Do they take an infinite number of (e) paths between states? |
| (2101) | Each of these models makes (eb) the same predictions, so one version is no more "true" than the others. |
| (2102) | What happens in practice is that we'll use whatever method is useful at the time. They are equivalent models, so the best model for the job is the one we'll use. The only reason the wavefunction view is so common is that it's (=) the version usually taught to introductory students. You might think this uncertainty of truth is due to (e) the behavior of quantum physics itself. It's so strange and counterintuitive that we can't (e) wrap our puny brains around what's really going on. But the same thing occurs in lots of other fields. Even something as straight forward as basic Newtonian physics |
| (2103) | Toss a baseball in the air and (eb) Earth's gravitational force pulls it down. |
| (2104) | The force of gravity is a simple truth, right? While we often describe classical motion in terms of (e&eb) forces and acceleration, we can also describe it in terms of (e&eb) energy and momentum. |
| (2105) | In the Lagrangian and Hamiltonian approach, the path of a baseball is (=) the optimized path among possibilities. |
| (2106) | In this view a baseball's path is (=) the extrema of an energy equation, and force can be derived as a necessary consequence of (e) this. |
| (2107) | In the relativistic view the baseball follows (e&eb) a geodesic, which is (=) the minimal path through (e&eb) space and time. |

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| (2108) So is a baseball's motion due to (e) a gravitational force, an energy extrema, or a spacetime geodesic? |
| (2109) As with quantum theory, different approaches yield (eb) the same result. They are mathematically equivalent, so we can use whatever method is most useful at the time. |
| (2110) At its core, science is less about truth and more about models. The metaphysics underlying a model is useful only as far as it allows us (eb) to make better predictions, generate (eb) new ideas, or bring models together as (=) a cohesive whole |
| (2111) This is why we have no problem using (e) classical gravity to calculate the path of (e) a spacecraft through (e&eb) the solar system, while (e) using special relativity to account for (e) the Doppler shift of the spacecraft's radio signals. |
| (2112) It's why we can use (e) quantum physics to study atoms in the morning, and general relativity to (e) study black holes in the afternoon. |
| (2113) In regimes where models conflict with each other it isn't a failure of truth, but instead shows an opportunity to develop a better model. |
| (2114) It could be that with each better model we move closer to (e&eb) the truth about reality. |
| (2115) The truth is out there, and science strives to move towards (e&eb) that truth. |
| (2116) It's a common view, and certainly the search for truth has driven (eb) many scientists to develop better and better models. |
| (2117) But the real power of science is the recognition that what we have are models. Our models can be powerful, but they are always a bit tentative. There's always a chance that they might just be an illusion of truth. |

**Synopsis: A Single-Photon Cheshire Cat July 1, 2016 Researchers detected the polarization of a photon separate from the photon itself, just as the grin of Lewis Carroll's Cheshire cat can appear apart from the cat's body Synopsis figure APS/Carin Cain**

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| (2118) In a quantum Cheshire cat, one property of a particle is present detached from (e) the particle itself, just as (=) the grin of the feline in Alice in Wonderland can exist in the absence of (e) the whole animal. |
| (2119) Now a research team has successfully separated (e&eb) a single photon from (e) its polarization. |
| (2120) The effect was previously demonstrated (eb) with neutrons and also with light in a classical scenario, but the latest experiment uses single photons. |
| (2121) To produce the Cheshire cat, **Maximilian Schlosshauer and his colleagues** at the University of Portland in Oregon used an interferometer where (e) light hits a beam splitter, goes in two, opposite directions around a circuit (but along slightly offset routes), and then recombines at (eb) the entrance. |
| (2122) The incoming polarization of light was horizontal, but this was switched to (e&eb) vertical for the counterclockwise loop. |
| (2123) Each path included a device to adjust (e&eb) the polarization further and another to (e&eb) scatter away a small number of photons of (e) a designated polarization, the equivalent of detecting their presence. |
| (2124) The team fired single photons into (e&eb) the interferometer so that only one photon was present at (eb) a time, traversing both paths simultaneously. |
| (2125) They placed a horizontal polarizer in front of the output photon detector and gradually varied (e&eb) the difference between the two path lengths. |

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| (2126) This variation generated (eb) a sine-wave interference pattern from (e) the exiting photons only when the polarization was tweaked in (eb) the counterclockwise path. |
| (2127) At the same time, the average number of photons counted decreased only when (e) light was scattered from the clockwise path. |
| (2128) Simultaneously demonstrating these different effects for two routes of a single photon constitutes (eb) evidence of the Cheshire cat. |
| **(2129)** The effect might someday improve precision experiments if (e) an unwanted property could be moved away from the property being measured. **This research is published in Physical Review A. – David Ehrenstein David Ehrenstein is the Focus Editor for Physics.** |

**Observation of the quantum paradox of separation of a single photon from one of its properties James M. Ashby, Peter D. Schwarz, and Maximilian Schlosshauer Phys. Rev. A 94, 012102 – Published 1 July 2016**

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| (2130) Authors report an experimental realization of the quantum paradox of (e) the separation of a single photon from (e) one of its properties (the so-called "quantum Cheshire cat"). |
| (2131) Authors use a modified Sagnac interferometer with (e&eb) displaced paths to produce (eb) appropriately pre- and postselected states of (e) heralded single photons. |
| (2132) Weak measurements of photon presence and (e&eb) circular polarization are performed in each arm of (e) the interferometer by introducing (eb) weak absorbers and small polarization rotations and analyzing changes in (eb) the postselected signal. |
| (2133) The absorber is found to have (e) an appreciable effect only in one arm of the interferometer, while (e) the polarization rotation significantly affects (e&eb) the signal only when (e) performed in the other arm. |
| (2134) Authors carry out both sequential and simultaneous weak measurements and find (eb) good agreement between measured and predicted weak values. |
| (2135) In the language of Aharonov et al. and in the sense of the ensemble averages described by (e) weak values, the experiment establishes (eb) the separation of a particle from one its properties during the passage through (e&eb) the interferometer. |

**Viewpoint: Relativity Gets Thorough Vetting from LIGO Frans Pretorius, Department of Physics, Princeton University, Princeton, NJ 08544, USA May 31, 2016• Physics 9, 52 The gravitational-wave signal observed by the LIGO detectors shows no deviation from what general relativity predicts. Figure caption expand figure Frans Pretorius, APS/Carin Cain**

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| (2136) Figure 1: The signal from one of the LIGO detectors in Hanford, Washington, is shown with two representations of (e) the best-fit numerical relativity (NR) waveform. |
| (2137) The filtered NR waveform illustrates how the raw waveform is perceived by (e) the detector, |
| (2138) On September 14, 2015, the LIGO interferometers made the first direct measurement of (e) gravitational waves [1]. |
| (2139) The historic discovery came nearly 100 years after Einstein showed that his general relativity theory predicts (eb) the existence of gravitational waves and over 50 years since (e) the first quest to observe the Universe in gravitational waves [2]. |
| (2140) The observed event, labeled GW150914, was (=) the collision of two black holes, one weighing 29 solar masses, the other 36 solar masses. |

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| (2141) This binary coalesced into (e&eb) a remnant black hole weighing 62 solar masses and rotating (e&eb) at 67% the maximum rate allowed (eb) for a black hole. |
| (2142) The most statistically significant portion of (e) the signal came from (e) the last several gravitational-wave cycles of the collision, where (e) the most nonlinear, strong-field dynamics of the theory manifest. |
| (2143) One couldn't have asked for a better first event to learn about (e&eb) this regime of general relativity. |
| (2144) The first thing that GW150914 can teach us is how closely nature follows (eb) Einstein's prescriptions for (e&eb) colliding black holes. |
| (2145) Researchers from the LIGO team and the Virgo team have performed a series of tests to see if (e) any violations of (e) general relativity might be lurking in the data of GW150914 [3]. |
| (2146) They first determined (eb) the theoretical waveform, or "template," from (e) general relativity simulations that best fit the measured signal (see Fig. 1). |
| (2147) Their main finding is that the residual signal—obtained by (e) subtracting this template from (e) the data—is consistent with (e&eb) noise. |
| (2148) This is, in a sense, the "everything" test, as (e) the template folds in (e&eb) all that general relativity predicts about the merger: the decay of the last few orbits of (e) the two black holes, their violent collision (e&eb) and astonishingly rapid relaxation to (e) a quiescent rotating black hole, the propagation of (e) gravitational waves through (e&eb) the cosmos, and finally the transformation of (e) these waves into(e&eb) a chirp in the LIGO detectors. |
| (2149) Given the large signal-to-noise ratio of 24, the close agreement between the template and (=) the data puts limits on (e) alternative models. |
| (2150) Specifically, the predictions of these models can only disagree with (e&eb) the general relativity prediction for GW150914 by, at most, 4%. |
| (2151) In addition to this comprehensive test, the LIGO/Virgo teams also performed other checks that pick apart the signal, in order to address particular aspects of general relativity and black hole physics. In one such test, they compared the early and (e&eb) late parts of the waveform. |
| (2152) From the waves generated during (e&eb) the initial inspiral, one can infer (eb) the masses and spins of the binary. |
| (2153) Given these numbers, general relativity allows (eb) the mass and spin of (e) the remnant black hole to be calculated. |
| (2154) The researchers compared this result to (e&eb) an independent measurement of (e) the mass and spin obtained from (e) the late-time ringdown, and to within (eb) uncertainties, the two estimates were (=) consistent. |
| (2155) Another test isolated coefficients in the so-called Post Newtonian (PN) series describing (eb) the inspiral. |
| (2156) The PN series is (=) an expansion of the gravitational-wave signal in powers of (e) the velocity vv of the sources relative to (e&eb) the speed of light cc. |
| (2157) Observations of binary pulsars, which show (eb) a decrease in their orbital period due to (e) gravitational-wave emission, have constrained (e) PN terms, especially at low order. |
| (2158) For higher-order terms associated with (e&eb) strong-field emission, the LIGO/Virgo teams were able to improve (eb) the limits by many orders of magnitude compared to (e&eb) the binary pulsar constraints. |

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| (2159) At first glance it might seem these new results simply add to (e) the litany of previous tests that general relativity passed with flying colors. |
| (2160) However, GW150914 is vastly more significant than that. Other tests have shown that general relativity is valid in (eb)the weak-field regime, but GW150914 gives (eb) the first direct observational evidence that general relativity is the correct theory of nature governing (e&eb) the extreme spacetime dynamics around the event horizon of (e) a black hole. |
| (2161) One could argue that binary pulsar observations already told us that gravitational waves exist [4], but GW150914 is more luminous in gravitational waves than pulsar binaries by (e) at least 24 orders of magnitude. |
| (2162) It's true that the evolution of the Universe as a whole is driven by (e) strong gravity, but it would be difficult to argue that precision cosmological measurements can constitute (eb) a test of general relativity when (e) one of the dominant inputs, dark energy, is still a mystery. |
| (2163) As far as black holes are concerned, astronomers have (e) observed ultra compact dark objects that are consistent with (e&eb) black holes. |
| (2164) The most striking data comes from (e) stellar orbits about the 4 million solar mass objects in (eb) the center of our Galaxy [5], but present observations cannot constrain it to be (=) a black hole as described by general relativity. |
| (2165) GW150914 has given us the first morsels of quantifiable evidence that the rich physics of black holes is as predicted by general relativity. For example, the LIGO observations are consistent with (e&eb) so-called cosmic censorship [6], which says that gravitational singularities must remain hidden behind an (eb) event horizon. |
| (2166) If a "naked" singularity were produced in (eb) the energetic collision of GW150914, the data would likely have shown (eb) some sort of anomaly, but in exactly what form, we don't know. |
| (2167) Another example involves the "no-hair" theorem [7], which says that the only properties of a black hole are (=) its mass, spin, and charge (though astrophysical black holes are expected to have negligible charge). |
| (2168) By consequence, these properties should determine (eb) the frequency and decay constants of (e) all the infinitely many ringdown modes of (e) a black hole. |
| (2169) With GW150914, only the least-damped ringdown mode is (=) visible above the noise, but when combined with (e&eb) the remnant mass and spin inferred from (e) the inspiral, this is enough to count as (=) the first piece of evidence for a black hole's baldness. |
| (2170) GW150914 should also be a test for (e) alternative theories of gravity or exotic alternatives to (e) black holes within general relativity, such as boson stars [8]. |
| (2171) Some of these alternative models are concerned with (e&eb) explaining dark energy or dark matter and do not (e) therefore, predict any deviation from relativity at (eb) the scale of GW150914. |
| (2172) However, other models do in fact deviate from (e) conventional relativity in their predictions of (e) the gravitational-wave signal from (e) a binary merger. |
| (2173) The largest deviations should presumably occur during (e&eb) the last several gravitational-wave cycles, where (e) the signal is strong, and deviations should be (=) easy to detect or constrain. |
| (2174) Embarrassingly though, the majority of alternatives that claim deviations can't (e) really say much about this event, as the two-body problem hasn't been solved in these theories. |
| (2175) These models can only reproduce (eb) a fraction of the full gravitational-wave signal, (eb) so the strong constraints one might have hoped for (e) cannot be leveled. |

(2176) Exotic alternatives within (eb) general relativity are in a similar boat, or worse for (e) objects like gravastars, which have (e) no well-defined theoretical framework to say what happens when (e) two such objects collide.

(2177) Regardless, what is clear from the data is that any non-black-hole solution must imagine a form of (e) exotic matter with a viscosity larger than (e) any known material to explain (eb) the rapid damping of oscillations observed during the ringdown.

(2178) The inferred event rate for heavy black hole mergers like GW150194 is high, and LIGO has another factor of 3 improvement in sensitivity slated before it reaches (eb) design sensitivity.

(2179) Within 5 years then, we can expect dozens of detected events, with the loudest having a signal-to-noise ratio upwards of (e) 100.

(2180) This will allow (eb) for exquisite insight into the nature of extreme gravity beyond (e) what GW150194 has already gifted us with. **This research is published in Physical Review Letters**.

**Tests of General Relativity with GW150914 B. P. Abbott et al. (LIGO Scientific and Virgo Collaborations) Phys. Rev. Lett 116, 221101 – Published 31 May 2016**

(2181) The LIGO detection of GW150914 provides (eb) an unprecedented opportunity to study the two-body motion of (e) a compact-object binary in (eb) the large-velocity, highly nonlinear regime, and to witness the final merger of (e) the binary and the excitation of uniquely relativistic modes of (e) the gravitational field.

(2182) Authors carry out several investigations to determine whether GW150914 is consistent with (e&eb) a binary black-hole merger in general relativity.

(2183) They find that the final remnant's mass and spin, as determined from (e) the low-frequency (inspiral) and (e&eb) high-frequency (postinspiral) phases of the signal, are mutually consistent with (e&eb) the binary black-hole solution in general relativity.

(2184) Furthermore, the data following the peak of GW150914 are consistent with (e&eb) the least-damped quasinormal mode inferred from (e) the mass and spin of the remnant black hole.

(2185) By using waveform models that allow (eb) for parametrized general-relativity violations during (e&eb) the inspiral and merger phases, authors perform quantitative tests on (e&eb) the gravitational-wave phase in the dynamical regime and they determine the first empirical bounds on (e&eb) several high-order post-Newtonian coefficients.

(2186) Authors constrain (e) the graviton Compton wavelength, assuming (eb) that gravitons are dispersed in (eb) vacuum in the same way as particles with mass, obtaining a 90%-confidence lower bound of 1013 km.

(2187) In conclusion, within our statistical uncertainties, we find no evidence for violations of (e) general relativity in the genuinely strong-field regime of gravity. **Received 26 March 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.221101 © 2016 American Physical Society**

**Focus: Background Noise of Gravitational Waves March 31, 2016• Physics 9, 33**

(2188) Assuming LIGO's recently detected black hole merger is not unusual; researchers revised upward their estimate of (e) the strength of the background noise coming from (e) distant mergers across the Universe. Figure caption expand figure Caltech/MIT/LIGO Lab (deleted due to spatial constraints)

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| (2189) Catching some waves The LIGO-Livingston interferometer consists of (e) two 4-km-long arms where (e) laser light bounces back and forth. Interference between the lasers from the two arms can signal tiny changes in (e&eb) the lengths of the arms resulting from gravi... |
| (2190) The recent detection of gravitational waves emitted from (e) a pair of merging black holes implies (eb) the existence of a large population of similar pairs, which should be generating (eb) a background of waves too faint to be detected with (e&eb) present technology. |
| (2191) Using this information, researchers now estimate that this background could be (=) roughly 10 times stronger than previously expected and might be found by (e) sensitive instruments in only a few years. |
| (2192) The background signal could help astrophysicists better understand how black hole binaries form and also pave the way to detecting the waves emitted by (e) the expanding Universe shortly after the big bang. |
| (2193) Gravitational waves were detected for the first time on 14 September 2015 by the Laser Interferometer Gravitational-Wave Observatory (LIGO). |
| (2194) By reflecting laser light between mirrors set 4 km apart, the team was able to detect (e&eb) tiny fluctuations in the distance between the mirrors as (=) the waves passed. |
| (2195) The waves were triggered by (e) an event, called GW150914; involving (e&eb) the merger of a pair of black holes having (e) masses of about 29 and 36 times that of the Sun. |
| (2196) Researchers expect that similar pairs of black holes, as well as neutron stars, should be merging constantly, creating (eb) gravitational waves that should bombard (e&eb) the Earth from all directions and make the equivalent of (=) a faint popping noise in a detector. |
| (2197) Previous estimates of the likely strength of this background had to make assumptions about (e&eb) how many pairs there might be and about the distribution of their masses. |
| (2198) These estimates remained speculative, as no actual binaries were known. |
| (2199) "What the recent observation has changed," says Nelson Christensen of Carleton College in Minnesota and the LIGO Scientific Collaboration (LSC), "is that we now know there really is (=) a population of black hole binaries. We have a small bit of information about it." |
| (2200) The new analysis—carried out by LSC researchers working with (e&eb) the European Virgo collaboration—estimates (eb) the energy density in the gravitational wave background in (eb) the frequency range below 100 Hz, where (e) the LIGO detector is most sensitive. |
| (2201) To make this estimate, they assumed (eb) a population of black hole binaries of various sizes, having (e) a range of distances from Earth, and created through (e&eb) a variety of scenarios. |
| (2202) They then calculated the amount of gravitational radiation that would result from (e) the merging of these objects. |
| (2203) The team then adjusted this estimate by using (e) details of the GW150914 event—especially the masses of the black holes involved—to make (eb) a most likely guess at the actual distribution of black hole binaries. |
| (2204) The idea is that the observed merger is (=) most likely representative of (e) the general population, rather than being extremely unusual. |
| (2205) If the black hole binary population is represented as (=) a bell curve, or Gaussian, the event should lie within (eb) the main bulge, not in the tails. |
| (2206) The observed mass of GW150914 suggests (eb) that the main population of black hole binaries could be (=) roughly 20 times more massive than expected, which would generate (eb) a background roughly 10 times "louder" than (e) previous estimates. |

(2207) However, Christensen emphasizes that the result has (e) lots of uncertainty, as it's based on (e) only one event.

(2208) "This estimate seems very plausible to me," says astrophysicist Maura McLaughlin of the University of West Virginia in Morgantown. She is a member of the NANOGRAV collaboration, which tries to detect gravitational waves at much lower frequencies than LIGO can observe. "I can't see any holes in their arguments. Unless LIGO got really lucky, then there must be many more of these systems than previously thought." This level of gravitational background implies (eb) that the LIGO and Virgo detectors might be sensitive enough to see it in 2–5 years, sooner than physicists previously expected, says Christensen, which has led to much anticipation. "Everyone is super excited," he says.

(2209) Detecting the background would provide (eb) a confirmation of the presence of gravitational waves carrying (e&eb) energy throughout the Universe. It would also be useful, the researchers suggest, in gaining (e) further understanding of precisely how black holes form.

(2210) And subtracting this background might someday allow (eb) detection of the fainter waves coming from (e) the expansion of the Universe moments after the big bang. This research is published in Physical Review Letters. –Mark Buchanan Mark Buchanan is a freelance science writer in Normandy, France.

**GW150914: Implications for the Stochastic Gravitational-Wave Background from Binary Black Holes B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 131102 – Published 31 March 2016**

(2211) The LIGO detection of the gravitational wave transient GW150914, from (e) the inspiral and merger of (e) two black holes with masses $\gtrsim 30 M_{\odot}$, suggests (eb) a population of binary black holes with (e&eb) relatively high mass.

(2212) This observation implies (eb) that the stochastic gravitational-wave background from binary black holes, created from (e) the incoherent superposition of (e) all the merging binaries in the Universe, could be (=) higher than previously expected.

(2213) Using the properties of GW150914, authors estimate the energy density of (e) such a background from (e) binary black holes.

(2214) In the most sensitive part of the Advanced LIGO and Advanced Virgo band for (e) stochastic backgrounds (near 25 Hz), authors predict $\Omega_{GW}(f=25\ \mathrm{Hz}) = 1.1^{+2.7}_{-0.9} \times 10^{-9}$ with 90% confidence.

(2215) This prediction is robustly demonstrated for (e) a variety of formation scenarios with (e&eb) different parameters.

(2216) The differences between models are (=) small compared to the statistical uncertainty arising from (e) the currently poorly constrained (e) local coalescence rate.

(2217) Authors conclude that this background is potentially measurable by (e) the Advanced LIGO and Advanced Virgo detectors operating at their projected final sensitivity. DOI:http://dx.doi.org/10.1103/PhysRevLett.116.131102 © 2016 American Physical Society

**Study opens new realms of light-matter interaction July 15, 2016 by David L. Chandler. Some sentences might have been omitted due to spatial constraints. Please bear with me.**

(2218) Emission spectra are a widely used method for (e) identifying chemical compounds; the bright lines reveal (eb) the different frequencies of light that can be emitted by (e) an atom.

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| (2219) Here, a normal emission spectrum for (e) an atom in a high-energy state (top) is compared to (e&eb) the emission from the same atom placed just a few nanometers (billionths of a meter) away from (e) graphene that has been doped with (e&eb) charge carriers (bottom). |
| (2220) For each energy-level transition, an orange line (or purple cloud) appears if (e) that transition is estimated to be faster than (e) one per microsecond — making it frequent enough to be (=) observed. **Credit: Massachusetts Institute of Technology** |
| (2221) A new MIT study could open up new areas of technology based on (e) types of light emission that had been thought to be (=) "forbidden," or at least so unlikely as to be (=) practically unattainable. |
| (2222) The new approach, the researchers say, could cause (eb) certain kinds of interactions between light and (e&eb) matter, which would normally take billions of years to happen, to take place instead (e&eb) within billionths of a second, under (e&eb) certain special conditions. The findings, based on a theoretical analysis, are reported today in the journal Science in a paper by MIT doctoral student Nicholas Rivera, Department of Physics Professor Marin Soljačić, Francis Wright Davis Professor of Physics John Joannopoulos, and postdocs Ido Kaminer and Bo Zhen. |
| (2223) Interactions between light and matter, described by (e) the laws of quantum electrodynamics, are the basis of a (e) wide range of technologies, including (e) lasers, LEDs, and atomic clocks. |
| (2224) But from a theoretical standpoint, "Most light-matter interaction processes are 'forbidden' by (e) electronic selection rules, which limits (e) the number of transitions between energy levels we have access to," Soljačić explains. |
| (2225) For example, spectrograms, which are used to (e) analyze the elemental composition of materials, show (eb) a few bright lines against (e&eb) a mostly dark background. |
| (2226) The bright lines represent (eb) the specific "allowed" energy level transitions in (eb) the atoms of that element that can be accompanied by (e) the release of a photon (a particle of light). |
| (2227) In the dark regions, which make up most of the spectrum, emission at those energy levels is (=) "forbidden." |
| (2228) With this new study, Kaminer says, "we demonstrate theoretically that these constraints can be lifted" using (e) confined waves within (e&eb) atomically thin, 2-D materials. |
| (2229) "We show that some of the transitions which normally take the age of the universe to happen could be made to happen within (eb) nanoseconds. |
| (2230) Because of this, many of the dark regions of a spectrogram become (=) bright once an atom is (=) placed near a 2-D material." |
| (2231) Electrons in an atom have (e) discrete energy levels, and when they hop from one level to another they give off (eb) a photon of light, a process called spontaneous emission. |
| (2232) But the atom itself is much smaller than the wavelength of the light that gets emitted—about 1/1,000 to 1/10,000 as big—substantially impairing (e&eb) the interactions between the two. |
| (2233) The trick is, in effect, to "shrink" (e) the light so it better matches the scale of the atom, as (=) the researchers show in their study. |
| (2234) The key to enabling a whole range of interactions, specifically transitions in (eb) atomic states that relate to absorbing or emitting light, is (=) the use of a two-dimensional material called graphene, in which light can interact with (e&eb) matter in the form of plasmons, a type of electromagnetic oscillation in the material. |
| (2235) These plasmons, which resemble photons but have wavelengths hundreds of times shorter, are (=) very narrowly confined in the graphene, in a way that makes some kinds of interactions with (e&eb) that matter many orders of magnitude more likely than (e) they would be in ordinary materials. |

(2236) This enables a variety of phenomena normally considered unattainable, such as (e) the simultaneous emission of multiple plasmons, or two-step light-emitting transitions between (e&eb) energy levels, the team says.

(2237) This method can enable the simultaneous emission of (e) two photons that are "entangled," meaning (eb) they share the same quantum state even when (e) separated.

(2238) Such generation of entangled photons is (=) an important element in quantum devices, such as those that might be used for cryptography.

(2239) Making use of these forbidden transitions could open up (eb) the ability to tailor the optical properties of materials in ways that had not been thought possible, Rivera says.

(2240) "By altering these rules" about the relationship between light and (e&eb) matter, "it can open (eb) new doors to reshaping (e&eb) the optical properties of materials."

(2241) Kaminer predicts that this work "will serve as (=) a founding piece for (e) the next generation of studies on light-matter interactions" and could lead to (eb) "further theoretical and experimental advances in many fields which rely on (e) light-matter interactions, including (e) atomic, molecular and optical physics, photonics, chemistry, optoelectronics, and many others."

(2242) Beyond its scientific implications, he says, "this study has possible applications across multiple disciplines, since (e) in principle it has potential to enable the full use of (e) the periodic table for optical applications."

(2243) This could potentially lead to (eb) applications in spectroscopy and sensing devices, ultrathin solar cells, and new kinds of materials to absorb solar energy, organic LEDs with higher efficiencies, and photon sources for (e) possible quantum computing devices."From the standpoint of fundamental science, this work lays the groundwork for a subfield that just a few years ago was difficult to imagine and until now was largely unexplored," Soljačić says.

(2244) "Two-dimensional materials confine (e) fields to a surface and motion to a plane, making possible (eb) many effects that are orders-of-magnitude too weak to (e) appear in a bulk volume," says Jason Fleischer, an associate professor of electrical engineering at Princeton University, who was not involved in this research.

(2245) This work, he says, "systematically explores how 2-D materials improve light-matter interactions, laying a theoretical foundation for faster electronic transitions, enhanced sensing, and better emission, including the compact generation of broadband and quantum light." More information: N. Rivera et al: Shrinking light to allow (eb) forbidden transitions on the atomic scale, Science (2016). DOI: 10.1126/science.aaf6308 Journal reference: Science search and more info website: Provided by: Massachusetts Institute of Technology search and more info website. This story is republished courtesy of MIT News (web.mit.edu/news office/), a popular site that covers news about MIT research, innovation and teaching.

**Viewpoint: Peering Back to the Beginning of Time Lawrence M. Krauss, Department of Physics and School of Earth and Space Exploration, Arizona State University, Tempe, AZ 85287, USA and Mount Stromlo Observatory, Research School of Astronomy and Astrophysics, Australian National University, Weston, ACT 2611, Australia June 19, 2014• Physics 7, 64**

(2246) The BICEP2 collaboration reports the detection of BB-mode polarization of (e) the cosmic microwave background—a signal that might originate from (e) gravitational waves created by (e) inflation during the very earliest moments in the evolution of our Universe. Figure caption expand figure P. A. R. Ade et al. [1]; Image on homepage: BICEP2 Collaboration (deleted due to restrictions on space)

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| (2247) Figure 1: (Top) BB-mode polarization pattern observed by (e) the BICEP2 telescope [1]. |
| (2248) The twisting pattern in the polarization of (e) the cosmic microwave background (CMB) might originate from (e) gravitational waves created by inflation in the very first moments of the universe. |
| (2249) Inflation—the hypothesis that the Universe underwent (e&eb) a phase of superluminal expansion in a brief period following the big bang—has (e) the potential of explaining, from first principles, why the Universe has (e) the structure we see today. |
| (2250) It could also solve outstanding puzzles of standard big-bang cosmology, such as why the Universe is, to a very good approximation, (=) flat and isotropic (i.e., it looks the same in all directions). |
| (2251) Yet we do not yet have a compelling model, based on (e) fundamental particle physics principles, that explains inflation. |
| (2252) And despite its explanatory power and a great deal of suggestive evidence, we still lack an unambiguous and direct probe of (e) inflation. |
| (2253) Theorists have developed different models for inflation, which all share (e&eb) a common, robust prediction: Inflation would have created (eb) a background of gravitational waves that could have (eb) an observable effect. |
| (2254) These waves would cause (eb) subtle, characteristic distortions of (e) the cosmic microwave background (CMB)—the oldest light in the Universe, released when (e) photons decoupled from matter and the Universe became (=) transparent to radiation |
| (2255) Until now, such a background of gravitational waves has (e) escaped detection. It is therefore very exciting that the BICEP2 (Background Imaging of Cosmic Extragalactic Polarization) collaboration has now reported in Physical Review Letters the observation of polarization modes of CMB radiation [1] that are consistent with (e&eb) such a background. |
| (2256) If confirmed, BICEP2's finding would provide (eb) the first direct probe of physical processes that occurred a mere 10−3610-36 seconds after (e) the big bang. |
| (2257) The result also offers a rare opportunity to directly test theoretical ideas, including inflation and grand unification (the theory that merges the electromagnetic, weak, and strong forces into a single force), at energies 13 orders of magnitude greater than (e) we can currently access at the Large Hadron Collider. |
| (2258) Inflation is based on (e) the possibility that a scalar field in the early Universe got momentarily stuck in (eb) a "false vacuum state"—a metastable condition in which the Universe wasn't at (eb) its global potential minimum. |
| (2259) The energy stored in this vacuum state could drive (eb) an exponential expansion of spacetime (similar to, but much faster than, the expansion due to dark energy seen today) before (e) a phase transition to the true vacuum state occurs, releasing (eb) this energy into (e&eb) particles and radiation. |
| (2260) This creates (eb) an initially hot state for (e) a subsequent big-bang-like expansion from (e) which our Universe subsequently evolved [2]. |
| (2261) If the Universe grew by at least a factor of 10261026, (eb) inflation can explain (eb) both the apparent flatness and isotropy of the Universe [2]: |
| (2262) The expansion would stretch (eb+ eb-dot) any initial curvature of (e) the Universe to near flatness, and smooth out (e&eb) any initial inhomogeneities, just as stretching (eb+ ebdot) a piece of fabric would do. |
| (2263) Isotropy is justified by the fact that the Universe, before inflation, was once much smaller and could have been contained in (eb) a causally connected region that was in thermal equilibrium, resulting in (eb) the approximately uniform, isotropic temperature distribution observed on (eb) large scales today. |

(2264) Importantly, the theory can explain the (eb) origin of structure in the Universe: Microscopic quantum fluctuations associated with (e&eb) the scalar field driving inflation could have been "frozen in" during the exponential expansion, resulting, (eb) after inflation, in density fluctuations that might have seeded all observed large-scale structures in the Universe, including (e) all visible galaxies and clusters of galaxies.

(2265) The signs of such primordial fluctuations have indeed been detected through (e&eb) their effect on (e&eb) the CMB, first by the Cosmic Background Explorer (COBE) [3] and later by other missions.

(2266) These observational studies detected small temperature variations (anisotropies) in (eb) the CMB and found a spectrum compatible with (e&eb) that predicted by inflation.

(2267) The magnitude of these fluctuations suggests (eb) an energy scale of inflation (loosely speaking, the energy stored throughout space in (eb) the scalar field potential that drove inflation) close to (e) 10161016 giga-electron-volts (GeV).

(2268) But while consistent with inflation, such fluctuations are not (e) conclusive evidence: Their predicted shape and magnitude is known to vary with (e&eb) the details of the inflationary potential (whose form is not yet experimentally constrained).

(2269) Hence many different observations could be accommodated, and any agreement between theory and observation is at best suggestive. There is also a possibility that other, noninflationary sources might contribute. However, gravitational waves—the propagating ripples in spacetime predicted by Einstein in 1916—might provide (eb) direct evidence of inflation.

(2270) During inflation, gravitational fields would also fluctuate because of (e) quantum effects.

(2271) These fluctuations would similarly get "frozen in" and re-emerge after inflation as (=) classical gravitational waves, whose amplitude depends only on (=) the energy stored in the field driving (eb) inflation, and not on other model-dependent details [4].

(2272) If the scale of inflation were indeed as high as a few times 10161016 GeV, this background could produce (eb) specific temperature anisotropies in the CMB [5].

(2273) These anisotropies have thus far not been detected by (e) the most recent searches by the WMAP and PLANCK satellites, and their absence was used to (e) put upper limits on the scale of inflation somewhat below 10161016 GeV.

(2274) But a more sensitive CMB probe of (e) gravitational waves exists.

(2275) Large-wavelength gravitational waves produced by (e) inflation would have alternatively compressed and expanded (e&eb) space, producing (eb) spatial variations in (eb) the intensity of radiation bombarding free electrons in (eb)the early Universe (see, e.g., Ref: [6]).

(2276) As photons decoupled from (e) matter and CMB radiation was released, scattering from (e) the undulating sea of electrons would have created (eb) patterns in the polarization of CMB.

(2277) Such polarization anisotropies are predicted in (eb) inflationary models to be (=) most visible when (e) looking at the sky on spatial scales comparable to (e&eb) or somewhat larger than the size of (e) the Universe's horizon at (eb) the time when the CMB was created.

(2278) This corresponds to (e&eb) about 1 degree of the field of view of the sky.

(2279) Over the past decade, various CMB polarization experiments have thus chased the observation of polarization anisotropies on (e&eb) these scales but, so far, also yielded (eb) merely upper limits on a possible gravitational background.

(2280) All of this has now changed with (e&eb) the BICEP2 collaboration's [1] report of (e) a signal that has the predicted (eb) characteristics of such a background at precisely the angular scales where (e) the inflationary signal is expected (see Fig. 1, top deleted due to spatial restraints).

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| (2281) The BICEP2 Collaboration uses (e) a microwave telescope located at the South Pole, with (e&eb) an array of polarization-sensitive detectors pointing at (eb) a small field (a solid angle of about 11 percent of the full sky) with a resolution of approximately 0.50.5 degrees, thus sufficient to probe the relevant angular scales for the detection of (e) polarization anisotropies originating from (e) **inflation**. |
| (2282) BICEP2's search is based on (e) two expected features of primordial gravitational waves. |
| (2283) First, gravitational waves would produce (eb) a characteristic spatial polarization pattern on CMB—a consequence of (e) their quadrupole nature. |
| (2284) One can decompose (e&eb) any spatial polarization pattern into (e&eb) two orthogonal components, a "curl-free" component and (e&eb) a component involving a twisting pattern. |
| (2285) In analogy to electromagnetic modes, these polarization patterns are called "EE modes" and (e&eb) "BB modes," respectively (see Fig. 1). |
| (2286) Scalar fluctuations in the density of (e) matter would produce (eb) EE modes, while (e) gravitational waves would instead generate (eb) a BB-mode pattern. |
| (2287) The second feature is the aforementioned characteristic spatial scale of (e) expected CMB polarization signatures, including (e) wavelength scales comparable to (e&eb) or larger than the horizon scale at the time the CMB was created. |
| (2288) Since inflation would produce (eb) both gravitational waves and (e&eb) scalar density perturbations, BICEP2, like all CMB-polarization experiments, quote their results in terms of (e&eb) the ratio rr of the amplitude of (e) tensor perturbations (associated with (e&eb) gravitational waves) to (e) the amplitude associated with (e&eb) scalar density fluctuations (reflected in (eb) the magnitude of CMB temperature fluctuations measured by (e) experiments like COBE). |
| (2289) A nonzero rr implies (eb) the detection of a BB-mode pattern and primarily measures (eb) the strength of any possible gravitational wave signal which, if (e) it comes from inflation, is then determined by (e) its energy scale (e.g., see Ref. [7] for a review). |
| (2290) BICEP2 reported a nonzero value of r=0.2(+0.07−0.05)r=0.2(+0.07-0.05), on scales ranging from (e) the horizon scale up to 4–5 times larger, with (e&eb) a claimed significance of better than 5σ5σ. |
| (2291) If confirmed, a value of 0.2 corresponds to (e&eb) an energy density during (e&eb) inflation of ∼1.2×1016∼1.2×1016 GeV. |
| (2292) The claimed signal is far larger than expected, partly because (e) the PLANCK satellite, despite not having yet reported its polarization measurements, has (e) previously put an upper limit on (eb) rr (<0.11<0.11). PLANCK's estimate is based on (e) the absence of a gravitational wave signature on (eb) the CMB temperature anisotropy on (eb) its largest scales (PLANCK observes the whole sky and therefore angular scales up to 360360 degrees). |
| (2293) The apparent discrepancy between the PLANCK upper bound and (e&eb) the BICEP2 detection might be a cause for concern. |
| (2294) However, the two bounds probe different spatial scales, and the BICEP collaboration has noted that, while tensor perturbations are expected to be (=) roughly scale independent, some scale dependence of (e) scalar density perturbations could relax (e&eb) the PLANCK constraint and bring (eb) the two results into agreement. |
| (2295) Of greater concern, of course, is the possibility that the signal could be determined by (e) foreground contamination—any source along (e&eb) the line of sight of the detector that could distort or contaminate (e&eb) the primordial CMB signal. |
| (2296) BICEP2 tested for (e) a wide variety of possible foregrounds, claiming (eb) they cannot be responsible for (e) the observed signal. |

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| (2297) One obvious candidate arises from (e) gravitational lensing: The intervening mass of (e) astrophysical bodies could distort (e&eb) the spatial pattern of polarization in a way that that could transform (e&eb) an EE-mode signal into (e&eb) a BB-mode signal. |
| (2298) The BICEP collaboration has analyzed this possibility extensively. Based on (e) measurements of (e) CMB temperature (and hence density perturbations) and on (e&eb) models of the mass distribution, they argue that such effects would only be (=) significant on smaller angular scales than those contributing to their measurements of rr. |
| (2299) But since BICEP2's first announcement on March 17, 2014, a number of scientists have raised another foreground concern due to (e) the possible impact of (e&eb) our Galaxy's dust. |
| (2300) BICEP2 modeled dust contamination scenarios based on (e) a variety of theoretical models, using data from (e) its predecessor BICEP1 and preliminary dust estimates from (e) the PLANCK satellite. |
| (2301) They concluded that dust could not (e) reproduce the magnitude of the observed signal. |
| (2302) However, PLANCK has in the meanwhile released (eb) more complete measurements of variation in polarized dust foregrounds in (eb) many regions of the Galaxy. |
| (2303) Their analysis increases the uncertainty associated with (e&eb) the dust contribution to (e) the BICEP2 result and also accommodates (e) a potentially larger polarized dust contribution. |
| (2304) Some researchers argued that dust polarization at a level 2–3 times higher than that assumed by BICEP2 is possible and could therefore plausibly be responsible for (e) the entire signal (e.g., see Ref: [8]). |
| (2305) While the BICEP2 collaboration continues to maintain that this is extremely unlikely based on (e) their estimates, many other researchers are more skeptical. |
| (2306) In any case, a variety of upcoming measurements, from PLANCK and other detectors are (=) expected to resolve this important ambiguity within the coming year. |
| (2307) Even if the BICEP2 signal does arise from (e) primordial gravitational waves, various consistency checks will be required (e) before it can be definitively argued that the signal arises from (e) inflation. |
| (2308) Inflation is certainly the best-motivated possibility, and the spectral shape of the signal, at least for (e) the five largest-scale modes observed by BICEP2, agrees extremely well with (=) inflationary predictions (see Fig. 1, bottom). |
| (2309) Several results have already appeared that strongly constrain (e) other possible sources, ruling out (e) for example, the possibility of gravitational waves generated by (e) other possible phase transitions after inflation [9] |
| (2310) . This conclusion is reinforced by more careful analyses of the scale dependence of (e) the predicted signal (e.g., Ref. [10]). |
| (2311) But it will nevertheless be important to assure that all possible alternative sources of (e) gravitational waves are ruled out. |
| (2312) If it can be established that the BICEP2 represents (eb) a gravitational wave signal from (e) inflation, there is one additional exciting implication. |
| (2313) Standard calculations suggest (eb) that the gravitational wave signal from inflation is due to (e) quantum fluctuations in the gravitational field, and this appears to require (e) the existence of gravitons—the hypothetical elementary particles that mediate (e&eb) the force of gravity. |
| (2314) Research by Frank Wilczek and myself, based on (e) a dimensional analysis argument [11], suggests that,(eb) independently of the method used to calculate the gravitational wave spectrum, a quantum gravitational origin is (=) required. |
| (2315) If so, the BICEP2 result then implies (eb) that gravity is ultimately a quantum theory, a result of fundamental importance for physics. |

(2316) It is hard to imagine a more important or wider-ranging cosmological observation. If confirmed by upcoming experiments, the BICEP2 results will push the frontiers of empirical physics forward in an unprecedented way, offering (eb) the first direct observation of gravitational waves and of quantum gravity effects as well as a new probe of the earliest moments in the Universe's life and of physics at grand-unified energy scales. This will thus constrain both our understanding of fundamental interactions and the evolution of our Universe as a whole.

**Detection of B-Mode Polarization at Degree Angular Scales by BICEP2 P. A. R. Ade et al. (BICEP2 Collaboration) Phys. Rev. Lett 112, 241101 – Published 19 June 2014**

(2317) Authors report results from the BICEP2 experiment, a cosmic microwave background (CMB) polarimeter specifically designed to search for (e) the signal of inflationary gravitational waves in (eb) the B-mode power spectrum around $\ell \sim 80$.

(2318) The telescope comprised a 26 cm aperture all-cold refracting optical system equipped with (e&eb) a focal plane of 512 antenna coupled transition edge sensor 150 GHz bolometers each with temperature sensitivity of $\approx 300\ \mu K_{CMB}\sqrt{s}$. BICEP2 observed from (e) the South Pole for three seasons from 2010 to 2012.

(2319) A low-foreground region of sky with (e&eb) an effective area of 380 square deg was observed to (e) a depth of 87 nK deg in Stokes Q and U.

(2320) In this paper authors describe the observations, data reduction, maps, simulations, and results. They find an excess of B-mode power over the base lensed-ΛCDM expectation in the range $30<\ell<150$, inconsistent with (e&eb) the null hypothesis at (eb) a significance of $>5\sigma$.

(2321) Through jackknife tests and simulations based on (e) detailed calibration measurements authors show that systematic contamination is (=) much smaller than the observed excess.

(2322) Cross correlating against WMAP 23 GHz maps we find that Galactic synchrotron makes a negligible contribution to the observed signal. We also examine a number of available models of polarized dust emission and find that at their default parameter values they predict power ~(5–10)× smaller than the observed excess signal (with no significant cross-correlation with our maps). However, these models are not sufficiently constrained by external public data to exclude the possibility of dust emission bright enough to explain the entire excess signal. Cross correlating BICEP2 against 100 GHz maps from the BICEP1 experiment, the excess signal is confirmed with $3\sigma$ significance and its spectral index is found to be consistent with that of the CMB, disfavoring dust at $1.7\sigma$. The observed B-mode power spectrum is well fit by a lensed-ΛCDM+tensor theoretical model with tensor-to-scalar ratio $r=0.20+0.07-0.05$, with $r=0$ disfavored at $7.0\sigma$. Accounting for the contribution of foreground, dust will shift this value downward by an amount which will be better constrained with upcoming data sets. Received 4 April 2014 DOI:http://dx.doi.org/10.1103/PhysRevLett.112.241101 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**(2323) Physicist claims to have observed quantum effects of Hawking radiation in the lab for the first time August 16, 2016 by Bob Yirka report**

(2324) This computer-simulated image shows a supermassive black hole at the core of a galaxy. The black region in the center represents the black hole's event horizon, where no light can escape the massive object's gravitational grip. The black …more

(2325) (Phys.org)—Jeff Steinhauer, a physicist at the Israel Institute of Technology, has published a paper in the journal Nature Physics describing experiments in which he attempted to create a virtual black hole in the lab in order to prove that Stephen Hawking's theory of radiation emanating from black holes is correct —though his experiments are based on sound, rather than light. In his paper, he claims to have observed the quantum effects of Hawking radiation in his lab as part of a virtual black hole—which, if proven to be true, will be the first time it has ever been achieved.

(2326) For many years, scientists believed that nothing could ever escape from a black hole. But in 1974, Stephen Hawking published a paper suggesting that something could—particles that are now called Hawking radiation. His idea was that if a particle (and its antimatter mate) appeared spontaneously at the edge of a black hole, one of the pair might be pulled into the black hole while the other escaped, taking some of the energy from the black hole with it—which would explain why black holes grow smaller and eventually disappear. Because such emissions are so feeble, no one has been able to measure Hawking radiation, so researchers have instead tried to build virtual black holes in labs to test the theory. One type of virtual black hole was proposed back in 1981 by Bill Unruh with the University of British Columbia—he suggested that an analogue might be created using water instead of light. He imagined a phonon existing at the edge of a waterfall—as the water speeds up, it begins to move faster than the speed of sound, causing it to be trapped. But if the phonon had an entangled mate that eluded the fall by moving away before getting caught up, it could escape. In this new effort, Steinhauer has built a device based on that idea and in so doing, claims he has observed an analogue of Hawking radiation.

(2327) The experiment consisted of creating an entangled pair of phonons sitting inside a bit of liquid that had been forced (via laser) to move very fast and then observing the action as one of the pair was pulled away as the liquid began to move faster than the speed of sound, while the other escaped—the fluid was a Bose-Einstein condensate of rubidium-87 atoms. After repeating the experiment 4,600 times Steinhauer became convinced that the particles were entangled, a necessity for a Hawking radiation analogue. His findings do not prove Hawking's theory to be true, of course, but they do appear to add a degree of credence that other researchers have thus far not been able to achieve. More information: Jeff Steinhauer. Observation of quantum Hawking radiation and its entanglement in an analogue black hole, Nature Physics (2016). DOI: 10.1038/nphys3863

**(2328) Abstract**

We observe spontaneous Hawking radiation, stimulated by quantum vacuum fluctuations, emanating from an analogue black hole in an atomic Bose–Einstein condensate. Correlations are observed between the Hawking particles outside the black hole and the partner particles inside. These correlations indicate an approximately thermal distribution of Hawking radiation. We find that the high-energy pairs are entangled, while the low-energy pairs are not, within the reasonable assumption that excitations with different frequencies are not correlated. The entanglement verifies the quantum nature of the Hawking radiation. The results are consistent with a driven oscillation experiment and a numerical simulation. Journal reference: Nature Physics search and more info website

Read more at: http://phys.org/news/2016-08-physicist-quantum-effects-hawking-lab.html#jCp

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as

the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t - T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

These equations can be written for

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(16) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.

(17) However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].

(18) Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.

Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

**TIME DELAY**

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period,

which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types system takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t - T_7)F_7(R_7(t - T_7) - N_7(t)M_7((R_7(t) \qquad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7}+\frac{N_7 dG_7}{dR_7}-\frac{dB_7}{dR_7}>\frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7}+\frac{dM_7}{dR_7}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also normalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_7$) and the other derived from the non-delayed or "immediate" processes($I_7$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_7=\frac{N_7 G_7}{2M_7}\frac{d(7nF_7)}{dR_7}(\geq 0) \qquad (7a)$$

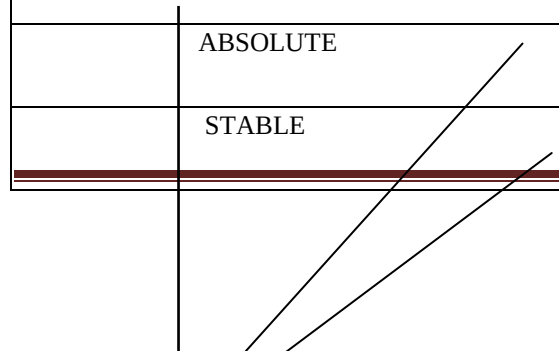
$$I_7=-\frac{N_7 G_7}{2M_7}\frac{d(7nM_7}{dR_7}(\geq 0) \qquad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7>L_7-I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

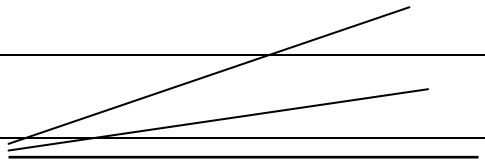
$$P_7>L_7-I_7, \quad \text{when} \quad P_7^2/4>I \qquad (8a)$$

$$P_7>L_7-I_7+\frac{(I_7-P_7^2}{(L_7-I)} \quad \text{when} \quad P_7^2/4<I. \qquad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,





**DELAYED PREDATOR CREDIT TYPES EFFECT WITH ITS WIDE RANGING RAMIFICATION**

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

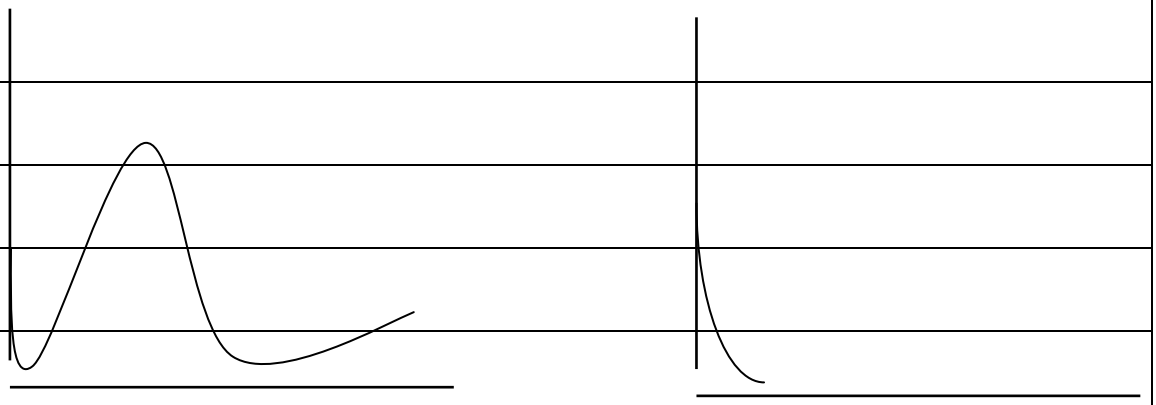
It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay

and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



**TIME DELAY**

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

To scores of problems this relationship is important. Of them some of the examples are as follows:

**New look at the galaxy-shaping power of black holes**

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| (641) | Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving. |
| (642) | "These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science. |
| (643) | "Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy." |
| (644) | This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return. |
| (645) | The remaining matter gets absorbed into the black hole, adding to (e) its mass. |
| (646) | The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne. |
| (647) | The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution. |
| (648) | Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form. |
| (649) | "Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara. |
| (650) | "It creates (eb) these bubbles, and the bubbles keep the plasma hot. |
| (651) | That's what prevents galaxies from (e) becoming even bigger than they are now." |
| (652) | Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space. |
| (653) | In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene. |
| (654) | Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve. |
| (655) | "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara. |
| (656) | "It's like nature's thermostat, which keeps these galaxies from (e) growing. |
| (657) | If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars." |

**A new look at the galaxy-shaping power of black holes**

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| (658) | An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the |
| (659) | McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said. |

(660) Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.

**Doomed satellite missions**

(661) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.

(662) At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.

(663) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.

(664) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.

(665) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.

(666) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people

(667) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."

(668) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.

(669) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(670) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.

(671) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.

(672) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.

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| (673) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses. |
| (674) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well. |
| (675) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma. |
| (676) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6. |
| (677) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus. |
| (678) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (679) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (680) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (681) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (682) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (683) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (684) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (685) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (686) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (687) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (688) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (689) We know dark matter exists only from (e) its gravitational effects. |
| (690) It holds together galaxies and (e&eb) galaxy clusters. |

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| (691) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (692) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |
| (693) They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (694) X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (695) They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (696) This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (697) An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |
| (698) Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (699) "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (700) The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (701) Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (702) Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (703) "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (704) Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (705) "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New |

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| York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (706) Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories. |
| (707) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen. |
| (708) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer. |
| (709) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants. |
| (710) These were treated as (=) "aware" and "unaware" trials, respectively. |
| (711) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it. |
| (712) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another. |
| (713) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion. |
| (714) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness. |
| (715) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said. |
| (716) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness." |
| (717) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner. |
| (718) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain. |
| (719) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said. |
| (720) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

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| (721) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter. |
| (722) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined. |
| (723) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions. |
| (724) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda. |
| (725) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general. |
| (726) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole. |
| (727) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes. |
| (728) Noda and his colleagues instead studied the change in emissions over (e&eb) time. |
| (729) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection. |
| (730) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1). |
| (731) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole. |
| (732) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved. |
| (733) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp |

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(734) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking “palm tree” organization.

(735) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society

(736) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible?

(737) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem.

(738) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other.

(739) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible.

(740) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges.

(741) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century.

(742) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching.

(743) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom.

(744) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth.

**Spheres on trees**

(745) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration.

(746) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement.

(747) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum.

(748) By visualizing the problem in this way, the researchers noticed that these particular graphs don’t (e) have lots of branches extending from lots of (e) other branches.

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| (749) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure. |
| (750) The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network. |
| (751) As a result, it provides (eb) an important clue for finding the global minimum. |
| (752) It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum. |
| (753) This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large. |

**Thompson's problem 2**

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| (754) By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society |
| (755) The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima. |
| **(756)** The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node. |

**Implications for Thomson's and Smale's problems**

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| (757) Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before. |
| (758) "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around." |
| (759) As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem. |
| (760) "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view." |
| (761) In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges. |
| (762) From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear. |
| (763) "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK. |
| (764) "For larger numbers, there will likely be (=) multiple funnels. |
| (765) Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum." |

(766) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.

(767) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.

(768) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t-T_7)F_7\big(R_7(t-T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t-T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) and the effective prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this

rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Look at the resemblance of String Theory and Abhinava's earlier theory and how the model could be applied: The world of manifestation and differentiated perceptions (Yikalpa) thus extends from one Centre to (eb) the next. **What is meant again is the realisation or the evolution of individual consciousness or general ledger to be (=)the cosmic general ledger is the sole criterion for the Satchidananda state where all dualities cease** Although it is never in fact divorced from (e) the subject who resides there, the ignorant fail (e) to grasp this fact and so, cut off from (e) the Centre, the world of objectivity becomes (=) for them the sphere of Maya. Bhagavatotpala quotes the Light of Consciousness amvitprakdsa): This ever pure experience (suddhanubhava) is variegated by (e) each form [revealed within it] Even so it remains unstained (nirmala) when moving to another. Just as a cloth which is naturally white, once dyed, cannot (e) change colour without [first] becoming white again, similarly the pure power of awareness, (citi) once coloured by (e) form, is pure [again] at the Centre where that form is (=) abandoned and from whence it proceeds to (e&eb) another. In his Essence of Vibration (Spandasarpdoha), K§emaraja explains that the rise and fall of every individual perception in the field of awareness is (=) a specific pulsation of consciousness. This is exactly the individual general ledger transaction lifted to the transaction in cosmic general ledger. During The Initial Instant Of Perception, T Consciousness Is Manifestly Apparent And The Yogi, Participating In Its Plenitude, Observes The Outer World Without Being Attached To Any Particular Divine Body And Sacred Circle Of The Senses Singling It Out From Any Other, Like A Man Who Observes A City From A High Mountain Peak.

He Sees The Outer World Reflected Within His Consciousness Free Of Thought-Constructs And So 'Stamps* The Outer On The Inner While Absorbing The Object And Means Of Knowledge In The Pure Subject Which Grasps Them As The Expansion Of His Own Nature. K§emaraja Says: By Penetrating Into Bhairavimudra, The Yogi Observes The Vast Totality Of Beings Rising From, And Dissolving Into, The Sky [Of Conscious- Ness], Like A Series Of Reflections Appearing And Disappearing Inside A Mirror. Through (e&eb) The Practice Of Bhairavimudra The Yogi Realises That He Is (=) The Substratum Consciousness (Adhisfhatr) Which Both Underlies (e&eb) And Is (=) The Essence Of All Things. He Discovers That Phenomena Have (e) No Independent Existence Apart From (e) Him And So Are, In This Sense, Void. At The Same Time, He Realises That Because (e) All Things Are Consciousness, They Are (=) Far From Unreal. He Views The Outer World Yet Sees (e) It Not. Beyond Both Voidness And Non-Voidness He Penetrates Into (e&eb) The Supreme Abode (Pararn Padam) Of Siva's Consciousness. [The Powers Of The Senses] Endowed With (e&eb) The Attributes Of The Great Union [Between Subject And (e&eb) Object] Whose Form Is The Wakening Of Man's Spiritual Potential (Kuntfalinf), Fill (e&eb) [With Consciousness] The Outer Clatter Of Diversity (Bhedatfambara) Born Of (e) Its Intense Power And Are Then Established In (eb) The Unobscured Abode Of The Void Of Consciousness To Shine [There] Eternally. Thus Residing Beyond (e) Being And Non-Being, The Sole Protector Of The Unity Which Is Tranquil And Expanding [Consciousness], Whose Glory Is All-Embracing And Form Unobscured, Is Called Bhairavimudra. Through (e&eb) The Practice Of Bhairavimudra, The Yogi Unites The Universal Vibration Of T Consciousness With (e&eb) The Individual Pulsation Of Objectivised 'This' Consciousness. The Two Aspects Of Consciousness Are Now In (eb) A State Of Equilibrium Like (=) The Two Pans Of An Evenly Weighed Balance And The Yogi Experiences (e&eb) The Pure Knowledge (Suddhavidya) That: 'I Alone Am (=)All Things'. Thus Becoming (=) The Master Of The Wheel Of Energies He Is (=) Free, Like Siva, To Create And Destroy. When [The Yogi] Is Well Established, Without (e) Wavering, Solely In The Integral Egoity Of His Authentic Nature, The Spanda Principle, And Is

Absorbed In Contemplation (Samavisfa), He Becomes (=) One With It The Doctrine Of Vibration (Tanmaya). Then . . . Dissolving And Creating The Universe By Means Of (er) His Introverted And Extroverted Absorption, He Destroys And Creates (e&eb) All Things Out Of Sarikara, His Innate Nature. [Thus] He Assumes The State Of The Universal Experiencer And Having Absorbed All That Is To Be Experienced From [The Grossest Level] — Earth — To (The Subtlest) — Siva — He Reaches (eb) The State Of The Supreme Subject By (e) Progressively Recognising [His Identity With Him]. Thus, Introverted And Extroverted Absorption Both Lead To (eb) The Recognition Of The Pulsation (Spanda) Of (e) One's Own Consciousness. This Is Where (e) Individual General Ledger Becomes (=) Cosmic General Ledger. At The Level Of Consciousness Corresponding To (e&eb) Siva's Basic State (Saryibhavdvastha), The Alternation From Inner To (e&eb) Outer Is Instantaneously Resolved Into (e&eb) The Vibration Of His Nature. When The Yogi Finally Comes To Be (=) Constantly Aware Of This Reality, His Enlightenment Is (=) Full And Perfect. Freed Of All Means (Anupaya) And Delighting In The Power Of His Bliss (Anandasakti), He Knows And Does Whatever He Pleases. The Yogi Seeking Self-Realisation Must Acquire (e) Mastery Over This Movement. Ksemaraja Stresses That The Doctrine Of Vibration Teaches That Liberation Can Only Be Achieved By (e) First Withdrawing All Sense Activity In Introverted Contemplation (Nimilanasamadhi) To (e) Then Experience The 'Great Expansion' (Mahdvikdsd) Of (e) Consciousness While (e) Recognising This To Be (=) A Spontaneous Process Within (eb) It. This Is Done Through (e&eb) The Practice Of Kramamudra. A Passage From The Now Lost Kramasiitra Explains: Although The Adept's Attention [May Be] Outwardly Directed, (eb) He Enjoys Contemplative Absorption Through (e&eb) The Introverted Aspect Of (e) Kramamudra. Initially He Turns Inward From (e)The Outside World And [Then] From Within [Himself] He Exits Into The Outer World Under The Influence Of His Absorption. There Is No Necessity Of Any Glorification Or Mortification Here. It Is Just Like Joining A Job And Retiring And Realising All That Happened Is Just By His Own Effort And Nothing Else. Thus The Sequence (Krama) In This Attitude (Mudra) [Ranges Through] Both Inner And Outer. The Yogi Must Pervade

The Surface Level Of Awareness (Yyutthana) With The Same Bliss He Experiences Plunging Into The Depths Of Con- Templative Absorption (Samadhi). Note Here That Samadhi Is A Death Like State And Still Being Matter And Information Stored You Experience Consciousness. Submerging Himself And Emerging Repeatedly From Samadhi, He Eventually Recognises That The Unity Of Consciousness Pervades Both States: The Best Of Yogi's, Who Has Achieved A State Of Complete Absorption Even When Risen From Meditation, [Inwardly] Vibrating Like A Drunkard In Blissful Inebriation From The After-Effects Of The Nectar Of Con- Templation, Sees All Things Dissolving In The Sky Of Consciousness Like A Cloud In The Autumn Sky. He Plunges Repeatedly within Himself Divine Body and Sacred Circle of the Senses and Becomes (=) Aware Of His Identity With (e&eb) Consciousness By (e) The Practice Of Introverted Contemplation. Thus Even When He Is Said To Have Risen From (e) Absorption, He Is (=) One With [His] Experience Of It. Ksemaraja Goes On To Explain That This Practice Is Called 'Mudra' Because (e) It Both Fills The Adept With (e&eb) Bliss (Mud) And Is Itself The Bliss Of (e) Consciousness. Moreover, It Dissolves (e) Away (Dra) All Bondage And 'Stamps' The Universe Of Experience With (e&eb)The Seal (Mudra) Of The Fourth State (Turiya) Of Enlightened Consciousness Beyond, And Including, The Three States Of Waking, Dreaming And Deep Sleep. It Is Called 'Krama' Because (e) It Is The Root Source Of (e) All Emanation And All Other Conscious Processes Which Succeed (e&eb) One Another In Ordered Sequence (Krama) And Is, (=)At The Same Time, Their Successive (Krama) Appearance As Well. By The Practice Of Kramamudra The Opposites Fuse And Siva And Sakti Unite. They Yogi Comes To Experience The Simultaneous Pervasion Of All The Lower, Grosser Categories Of Existence By The Higher And The Presence Of The Lower In The Higher. Commencing His Practice In A Low Form Of Bhairavimudra, The Yogi Conjoins The Outer With The Inner; Then, In Kramamudra, He Fills Both The Outer With (e&eb) The Inner And (e&eb) The Inner With (e&eb)The Outer. When He Achieves (eb) Perfection In This Two-Fold Movement, He Attains To (e) The Highest Form Of Bhairavimudra In Which The Two Merge Completely In (eb)The Experience Of (e) The Absolute (Anuttara), Free

Of (e) All Differentiation And (e&eb) Polarities. If He Fails To Maintain Awareness Of This State, He Again Falls Into (eb) Kramamudra Until He Has Finally Completely Merged All The Highest States In (e&eb) The Lower And The Lower In (e&eb) The Higher. He Then No Longer Needs To (e) Resort To (e) Any Means (Anupaya) To (e) Achieve Liberation. All He Says or Does Anything He Perceives or Thinks, Instantly Occasions In (eb) Him The Highest Level Of Consciousness. Thus The Fruit Of Bhairavimudra Is (*=)The Wonder (Camatkara) Or Amazement (Vismaya) That Overcomes (e)The Yogi When (e) He Reaches The Plane Of (e) Union (Yogabhumika) Y Xsl Where (e) All Opposites Merge In (e&eb)The Radiance Of The Great Light Of Consciousness. The Stanzas On Vibration Teach: How Can One Who, As If Astonished, Beholds (e) His Own Nature As (=) That Which Sustains [The Existence Of Everything] Be Subject To (e&eb) This Painful Round Of Transmigration (Kusrti)? The Yogi, Recognising His True Nature To Be (=)The Supreme Subject, Is Astonished To (e) Suddenly Discover That The Individual He Thought He Was, Caught Up (e&eb) In The Trammels Of Thought And Living In A World Enmeshed In (e&eb) The Web Of Time And Space, Does Not Really Exist At All. He Experiences A 'Turning About' (Paravrtti) In (eb)The Deepest Seat Of The Doctrine Of Vibration Consciousness As He Penetrates His True Nature. The Sudden Eruption Of This Intuition (Pratibha) Arouses In Him A Cry Of Amazement As He Transcends All Thought-Constructs And, Perfectly Absorbed In His Own Nature, Is Liberated. The Path To Liberation Essentially, Spanda Doctrine Is Concerned With Two Matters. The First Is To Impart To Those Who Are Fit To Receive The Teachings A Deeper Understanding Of The Ultimate Goal Of Life (Upeya). When We Have Understood What Truly Benefits Us And Is Worth Attaining And What, On The Contrary, Is Of No Real Value But Stands In The Way Of This Attainment, We Can Begin To Make Progress Towards Our Goal. This Is Spanda Doctrine's Second Concern, Namely, To Show The Way In Which We Can Develop Spiritually Through Siva's Grace And The Right Application Of The Means To Realisation That It Teaches. When Both These Aspects Of The Teaching Have Been Correctly Understood And Applied, The Spanda Yogi Achieves A Clear And Permanent

Realisation Of His Goal And Is Liberated, Thus Fulfilling The Ultimate Aim Of The Teaching. The Doctrine Of Vibration Is Not Meant For The Spiritually Dull. It Is Not For The Worldly Whose Consciousness, Clouded By Ignorance, Is As If Dreaming, Even During The Waking State Of Daily Life, The Dream Of Its Own Thought-Constructs. The Teachings Are Meant For Those Who Are Awake (Prabuddha), Those Who, Full Of Faith And Reverence, Are Always Alert And Intent On Discerning The True Nature Of Ultimate Reality. This Reality Is Understood In Three Basic Ways. The First Is Purely Transcendental. The Stanzas Choose This Aspect As The One Which Formally Defines It Most Specifically. Ultimate Reality Transcends All The Opposites, Including Subject And Object. This Does Not Mean, However, That It Is An Unconscious Void, A Mere Absence Of All Existence. In Fact, This Negative Characterization Of Reality (Which Includes Also A Denial Of All That Is Unconscious) Implies A Positive Immanence In Which The Opposites Are United In The Oneness Of Pure Consciousness That Is Equally Siva And Spanda, His Universal Activity. These Two Seemingly Contrasting Aspects Are Reconciled In The Third, Namely, Reality Understood As The Essential Nature Of All Things. Although Universal And Everywhere The Same, It Is Understood To Be The Essential And Specific Nature Of Each Existent As Its 'Own Nature' (Svabhava).

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such

a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

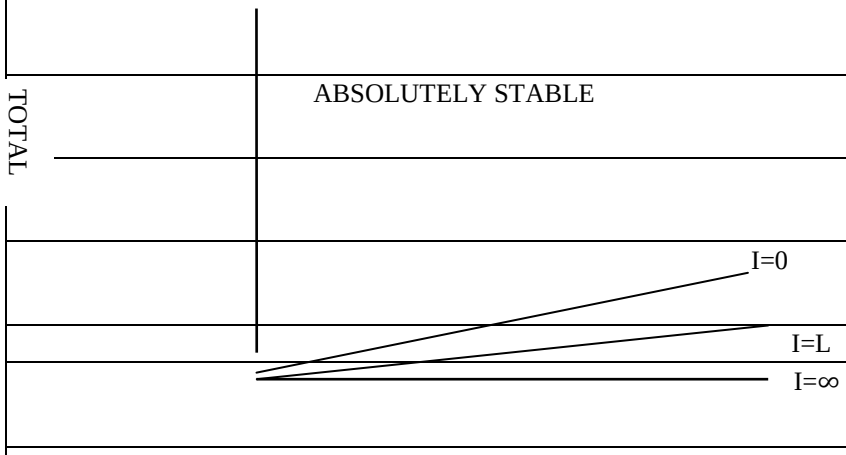
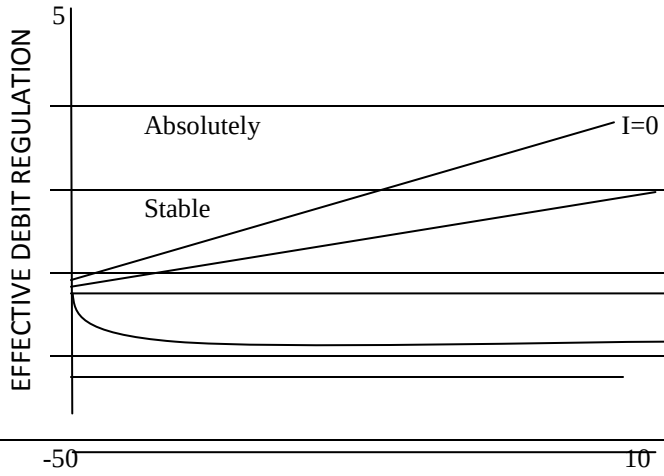



FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

This equation is of paramount importance in the following cases: According to the principles of geometric quantum mechanics, the physical characteristics of a given quantum system can be represented by geometrical features that are preferentially identified in this complex manifold. Here we construct a number of examples of such features as they arise in the state spaces for spin, spin 1, spin and spin 2 systems, and for pairs of spin systems. A study is then undertaken on the geometry of entangled states. A locally invariant measure is assigned to the degree of entanglement of a given state for a general multi-particle system, and the properties of this measure are analysed for the entangled states of a pair of spin particles. With the specification of a quantum Hamiltonian, the resulting Schrödinger trajectories induce an isometry of the Fubini–Study manifold, and hence also an isometry of each of the energy surfaces generated by level values of the expectation of the Hamiltonian. For a generic quantum evolution,

the corresponding Killing trajectory is quasiergodic on a toroidal subspace of the energy surface through the initial state. When a dynamical trajectory is lifted orthogonally to Hilbert space, it induces a geometric phase shift on the wave function. The uncertainty of an observable in a given state is the length of the gradient vector of the level surface of the expectation of the observable in that state; a fact that allows us to calculate higher orders corrections to the Heisenberg relations. A general mixed state is determined by a probability density function on the state space, for which the associated first moment is the density matrix. The advantage of a general state is in its applicability in various attempts to go beyond the standard quantum theory, some of which admit a natural phase-space characterization. **Journal of Geometry and Physics Volume 38, Issue 1, April 2001, Pages 19–53 Geometric quantum mechanics Dorje C. Brody Lane P. Hughston DOI: 10.1016/S0393-0440(00)00052-8** Expansion (unmesa) of [consciousness] or the creative intuition {pratibha) is (=) [experienced] in the interval which divides two [moments] of differentiated perception (vikalpa) It is here that they arise and disappear. The Sastras and Agamas proclaim with reasoned argument that it is (=) free of thought-constructs {Nirvikalpa) and precedes (e) Path to Liberation all mental representations of any object. None can deny that a gap exists between perceptions insofar as two moments of thought are (=) invariably divided. This [gap] is (=) the undifferentiated unity of all the countless manifestations. Similarly, in the outer more objective sphere, where change consists (e) of the alterations in the configurations of manifest appearances (Abhasa\ the transition from one to (e&eb) another corresponds to (e&eb) a phase of pure luminosity that marks (eb) the beginning of one form and the end of another

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^{\infty} \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7 \left( U_7 - V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7 \right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$then given the constraints of $f(T_7)$, $(A1.5) can\ never\ be\ satisfied\ Co$nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7 \lambda_7 e^{-T_7 \lambda_7} + [(U_7 - V_7) M_7 + Y_7] - [((U_7 -$$

V 7M 7−]**e**−7 7λ 7=0

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$\boldsymbol{Z}_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7) \cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7) \sin(Z_7 T_7) = 0$$

The warrant of Truth is ever elusive when we deconstruct the foundational justifications of our convenient notions about the way the world works. It is easy to confuse what is actually the creation of beliefs with the "discovery of Truth," a common goal of science and theology. Science offers no ontological Scientific Picture of the World but substitutes a number of applied theories as meta-theory. It's close, powerful and useful: it 'works' and allows use to do work. But the still mysterious "true" metatheory may reveal the basis of consciousness. But who knows when we will lift that veil? Permeating the living reality of our culture are certain contagious notions, fads and trends that have the ability to influence the way we think about the nature of Reality and ourselves. Some of them are toxic and can consume you; others are just whims of pop culture infectious and benign as the common cold. **Jung** described the concept of psychic contagion by certain archetypal forces inherent in the human psyche, which manifest in our spiritual lives and belief systems. He spoke of both conscious and unconscious contamination. Notions like this range from simple superstitions to scientific concepts, to urban myths. Notions sweep through our culture and insinuate themselves within its fabric, as fads, whether they are "real" or not, they can be influential. They mood alter us, make us feel we belong. An analysis of these notions is useful in distinguishing a common human phenomenon from any potential "alien influence" which may or may not be exerted on us from an unknown source. **Strange Attractors: Transference, Holography, and an Archetype Burke, J. (2003).**

**Strange Attractors: Transference, Holography, and an Archetype (Doctoral dissertation, Pacifica Graduate Institute, 2003).** In consideration to the Nobel Proze for Chemistry this year(2014), a stability nalysis of the proble in to single molexues would be a contemporaneous question that has to be taken cognisace of and given importance to . Determination of molecular structure by geometry optimization became routine only after efficient methods for calculating the first derivatives of the energy with respect to all atomic coordinates became available. Evaluation of the related second derivatives allows the prediction of vibrational frequencies if harmonic motion is estimated. More importantly, it allows for the characterization of stationary points. The frequencies are related to the eigenvalues of the Hessian matrix, which contains second derivatives. If the eigenvalues are all positive, then the frequencies are all real and the stationary point is a local minimum. If one eigenvalue is negative (i.e., an imaginary frequency), then the stationary point is a transition structure. If more than one eigenvalue is negative, then the stationary point is a more complex one, and is usually of little interest. When one of these is found, it is necessary to move the search away from it if the experimenter is looking solely for local minima and transition structures. The total energy is determined by approximate solutions of the time-dependent Schrödinger equation, usually with no relativistic terms included, and by making use of the Born–Oppenheimer approximation, which allows for the separation of electronic and nuclear motions, thereby simplifying the Schrödinger equation. This leads to the evaluation of the total energy as a sum of the electronic energy at fixed nuclei positions and the repulsion energy of the nuclei. Notable exceptions are certain approaches called direct quantum chemistry, which treat electrons and nuclei on a common footing. Density functional methods and semi-empirical methods are variants on the major theme. **Orlando Alvarez** discusses Polyakov's quantization of the string in the presence of a boundary allowing for an arbitrary topology for the world sheet. In addition to the dynamical conformal factor discovered by Polyakov, there are a finite

number of new degrees of freedom if the surface is more complicated than a sphere or a disc. The quantization of the Liouville theory in an arbitrary topology is also discussed. A one-loop calculation shows that the model is renormalizable if one performs a mass renormalization and an additive field renormalization. The renormalization group equations have a perturbative infrared unstable fixed point in all topologies. Copyright © 1983 Published by Elsevier B.V. **Nuclear Physics B Volume 216, Issue 1, 25 April 1983, Pages 125–184 Theory of strings with boundaries: Fluctuations, topology and quantum geometry Orlando Alvarez DOI: 10.1016/0550-3213(83)90490-X** Quantum geometry (the modern loop quantum gravity involving graphs and spin-networks instead of the loops) provides microscopic degrees of freedom that account for black-hole entropy. However, the procedure for state counting used in the literature contains an error and the number of the relevant horizon states is underestimated. In our paper a correct method of counting is presented by **Marcin Domagala and Jerzy Lewandowski.** Results lead to a revision of the literature of the subject. It turns out that the contribution of spins greater than 1/2 to the entropy is not negligible. Hence, the value of the Barbero–Immirzi parameter involved in the spectra of all the geometric and physical operators in this theory is different than previously derived. Also, the conjectured relation between quantum geometry and the black-hole quasi-normal modes should be understood again. **Marcin Domagala and Jerzy Lewandowski 2004 Class. Quantum Grav 21 5233 doi:10.1088/0264-9381/21/22/014 Black-hole entropy from quantum geometry** Quantum geometry predicts that a universe evolves through an inflationary phase at small volume before exiting gracefully into a standard Friedmann phase. This does not require the introduction of additional matter fields with ad hoc potentials; rather, it occurs because of a quantum gravity modification of the kinetic part of ordinary matter Hamiltonians. Authors draw the cognizance of application of the same mechanism that can explain why the present day cosmological acceleration is so tiny. DOI: http://dx.doi.org/10.1103/PhysRevLett.89.261301

**Inflation from Quantum Geometry Phys. Rev. Lett 89, 261301 – Published 12 December 2002 Martin Bojowald** The loop quantum cosmology of the closed isotropic model is studied with special emphasis on a comparison with traditional results obtained in the Wheeler-DeWitt approach by **Martin Bojowald** .Investigation includes the relation of the dynamical initial conditions to boundary conditions such as the no-boundary or the tunneling proposal and a discussion of inflation from quantum cosmology. **DOI: http://dx.doi.org/10.1103/PhysRevD.67.124023 Loop quantum cosmology, boundary proposals, and inflation Phys. Rev. D 67, 124023 – Published 23 June 2003 Martin Bojowald and Kevin Vandersloot** In loop quantum cosmology, the universe avoids a big bang singularity and undergoes an early and short super-inflation phase. During super-inflation, non-perturbative quantum corrections to the dynamics drive an inflaton field up its potential hill, thus setting the initial conditions for standard inflation. **Shinji Tsujikawa et al** show that this effect can raise the inflaton high enough to achieve sufficient e-foldings in the standard inflation era. They also analyse the cosmological perturbations generated when slow-roll is **violated after super-inflation** and show that loop quantum effects can in principle leave an **indirect signature** on the largest scales in the CMB, with some loss of power and running of the spectral index. **Shinji Tsujikawa et al 2004 Class. Quantum Grav 21 5767 doi:10.1088/0264-9381/21/24/006 Loop quantum gravity effects on inflation and the CMB Quantum geometry from quantum information** is the study of **Florian Girelli and Etera R Livine .**Loop quantum gravity defines the quantum states of space geometry as spin networks and describes their evolution in time. We reformulate spin networks in terms of harmonic oscillators and show how the holographic degrees of freedom of the theory are described as matrix models. This allows us to make a link with non-commutative geometry and to look at the issue of the semi-classical limit of loop quantum gravity from a new perspective. This work is thought of as part of a bigger project of describing quantum geometry in quantum information terms.

**Florian Girelli and Etera R Livine 2005 Class. Quantum Grav 22 3295 doi:10.1088/0264-9381/22/16/011 Reconstructing quantum geometry from quantum information: spin networks as harmonic oscillators.** Quantum geometric mechanics is the topic of study by **Dorje C. Brody Lane P. Hughston** establishing a clear relation between the quantum mechanics and quantum geometry and also the entangled states. The manifold of pure quantum states can be regarded as a complex projective space endowed with the unitary-invariant Fubini–Study metric. Note how differential gheometry concepts with Nunney's model could be entangle so that it is feasibly and viably applied to various type of model including consciousness (See Doctrine of Vibration)

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $\boldsymbol{Z}_7^2$ and $Z_7$in (A2.2):

$W_7^2 (X_7^2$-3 $M_7\ A_7\ X_7$- $M_7^2 (X_7$- $Y_7) + 2\ M_7^3\ A_7$+2 $A_7^2\ M_7^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_7$=1-cos ($Z_7 T_7$).

In the absence of a delay the system is stable for $A_7 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7 > (X_7 - Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$A_7 >$ [ $X_7 - Y_7/2\ M_7$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_7, A_7) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_7 \le 2$ *is* of interest First consider the extreme values

$$Q_7(0, A_7) = A_7^2 (X_7 - Y_7)$$

$$Q_7(2, A_7) = (A_7 + 2M_7)^2 (2A_7 M_7 - (X_7 - Y_7))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7 M_7^2 - 3M_7 X_7\right) + W_7\left(3M_7 A_7^2 - 2A_7 X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7 + Y_7\right)$$

$$=2M_7^2 W_7^2 (M_7 + A_7) + M_7 A_7^2 W_7 + (2A_7 M_7 - X_7 + Y_7)(A_7 W_7 + M_7 W_7^2$$

$$+ (2 - W_7)((2M_7 X_7 + M_7 Y_7) W_7 + A_7 (X_7 + Y_7)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $W_7$touches the $W$–axis only at its minimum value in which case (A2.9) shows that

further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$has no minimum in the range 0≤ $W_7$ ≤2.Defining

$Q_7(W_7, A_7) = a_7 W_7^2 + c_7$ *where* $a_7\, b_7$ *and* $c_7$ Are functions of $A_7$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7 \geq b_7 \geq 0$$

And

$$b_7^2 \geq 4a_7 c_7$$

$$Case(i): Y_7 = 0$$

*Let* $Y_7 = 0$ And try the substitution A= $X_7 = /2M_7\, in$ (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_7 X_7 /$ $2M_7$) is positive for *W<2. Since* $Q_7$ increases with $A_7$(see A2.9), increasing $A_7$will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

$Case\ (ii): Y_7 = y_7\, X_7$ ,

$Let\ Y_7 = y_7\ X_7$And try the substitution $A_7 = X_7/(2 + A_7\ y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + n_7 y_7)$ we have

$$W_7^2\ y_7\left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2\left(2 - n_7 + n_7 y_7\right)\right]$$

$$-W_7\ X_7\left[\frac{X_7(1 + n_7 y_7)}{M_7^2(2 + n_7 y_7)^2} + 2\ y_7\left(n_7 - 1\right) + \frac{X_7^2(7 + y_7)}{M_7^2(2 + n_7 y_7)} = 0.\ (A2.12)\right.$$

Since $n_7$ and $A_7$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7), the\ condition (A2.77?) can\ be\ written\ as$

$[\frac{X_7(E_7-7)}{E_7^2}]^2[\frac{X_7^2}{M_7^2} - 4\ y_7\ E_7^2\frac{X_7}{M_7^2} + 4y_7(n_7 - 2)E_7^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$2\ E_7[y_7 E_7 + (y_7\ E_7(7 + y_7)(E_7 - n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - y_7 E_7(7 + y_7)(E_7 - n_7))^{7/2}]$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0 $W_7$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7-7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2-2y_7\frac{(E_7+n_7^2y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right)-4y_7(E_7-n_7)$$

And has the solution

$$\frac{X_7}{M_7^2}>\frac{E_7}{(E_7-7)}[y_7(E_7+n_7^2y_7)+[y_7^2(E_7+n_7^2y_7)^2+4y_7(E_7-7)(E_7-n_7)]^{7/2}]$$

Given $E_7>n_7$this lower limit to $X_7/M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7/M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$in the appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7<n_7$ $(ie.,2A_7M_7^2<X_7-Y_7$)then (A2.7)and (A2.8)show that $Q_7$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2}>2E_7[y_7E_7-(y_7E_7(7+y_7)(E_7-n_7)]^{7/2}]$$

And

$$E_7>n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2 > 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.

If $A_7^2 < 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.+ $(Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$ and $X_7$ and $Y_7$ by $W_7^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7 \lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]\, e^{2T_7\lambda_7} + M_7 V_7\, e^{2T_7\lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i} Z_7$ *and* separate

(A3.1)into real and imaginary parts

$$Z_7^2 + (M_7 + V_7) Z_7 \sin(Z_7 T_7) - (U_7 M_7 + Y_7) + (M_7(U_7 + V_7) - X_7)$$

$$\mathrm{X} \cos(Z_7 T_7) - 2M_7 V_7 \cos^2(Z_7 T_7) + M_7 V_7 = 0$$

$$(M_7 + U_7) Z_7 - (M_7 + V_7) Z_7 \cos(Z_7 T_7) + (M_7(U_7 + V_7) - X)$$

$$\mathrm{X} \sin((Z_7 T_7) - M_7 V_7 \sin(Z_7 T_7) = 0$$

Using (A.3.3) to eliminate $Z_7$ and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - cos\, Z_7\ T_7$ and $A_7 = U_7 - V_7$ we obtain

$$4\,M_7^2 V_7^2 W_7^4 + 2\,M_7 V_7\ W_7^3[(M_7 - V_7)^2 + A_7(M_7 + V_7) + 2(A_7 M_7 - X_7)]$$

$$+ W_7^2[M_7 - V_7)^2(2\,A_7\ M_7 + Y_7 - X_7) + (2A_7\ M_7 - X_7)^2 - A_7(M_7 - V_7)(2A_7\ M_7 - X_7)$$

$$+ M_7 V_7\,(A_7^2 + 4(X_7 + Y_7))]$$

$$+ W_7[A_7^2 + 2X_7)(2\,A_7\ M_7 - X_7) + 2A_7 Y_7(M_7 + V_7) - M_7 A_7^3] + A_7^2(X_7\ M_7 - Y_7) = 0$$

Being a quartic in $W_7$ (A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7 > 0$$

$$and$$

$$C_7 > (X_7 - Y_7)/2\,M_7$$

*recalling that* $A_7 = U_7$ and defining $C_7 = U_7 + V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0, U_7 V_7 = A_7^2(X_7 + Y_7)$$

$$Q_7\left(2, U_7 V_7\right) = \left(C_7 + 2M_7^2\right)^2 (2C_7 M_7 + Y_7 - X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

## Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

**Module numbered five**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This is the condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference

between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \quad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t - T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\left(R(t)\right)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

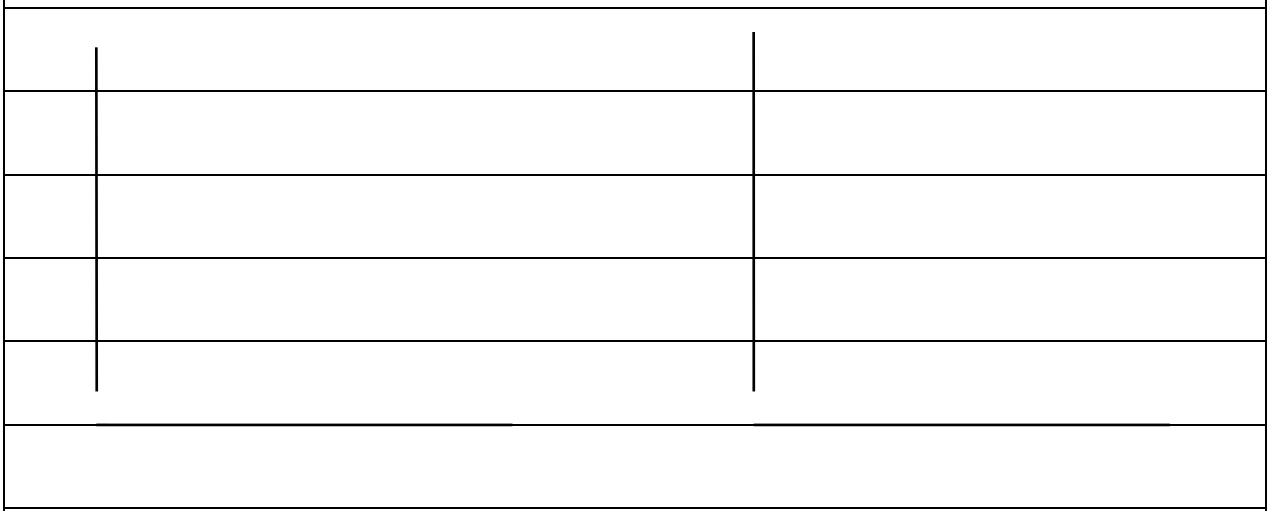
Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time

delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the



TIME DELAY

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period,

which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

As it is usually expressed outside of phenomenological texts, to undergo a conscious experience necessarily means that there is something it is like for the subject to have that experience (Nagel 1974; Searle 1992). This is obviously true of bodily sensations like pain. But it is also the case for perceptual experiences, experiences of desiring, feeling, and thinking. There is something it is like to taste chocolate, and this is different from what it is like to remember what it is like to taste chocolate, or to smell vanilla, to run, to stand still, to feel envious, nervous, depressed or happy, or to entertain an abstract belief. Yet, at the same time, as I live through these differences, there is something experiential that is, in some sense, the same, namely, their distinct first-personal character. All the experiences are characterized by a quality of mineness or for-me-ness, the fact that it is I who am having these experiences. All the experiences are given (at least tacitly) as my experiences, as experiences I am undergoing or living through. All of

this suggests that first-person experience presents me with an immediate and non-observational access to myself, and that consequently (phenomenal) consciousness consequently entails a (minimal) form of self-consciousness. To put it differently, unless a mental process is pre-reflectively self-conscious there will be nothing it is like to undergo the process, and it therefore cannot be a phenomenally conscious process. The mineness in question is not a quality like being scarlet, sour or soft. It doesn't refer to a specific experiential content, to a specific what; nor does it refer to the diachronic or synchronic sum of such content, or to some other relation that might obtain between the contents in question. Rather, it refers to the distinct givenness or the how it feels of experience. It refers to the first-personal presence or character of experience. It refers to the fact that the experiences I am living through are given differently (but not necessarily better) to me than to anybody else. It could consequently be claimed that anybody who denies the for-me-ness of experience simply fails to recognize an essential constitutive aspect of experience. Such a denial would be tantamount to a denial of the first-person perspective. It would entail the view that my own mind is either not given to me at all — I would be mind- or self-blind — or is presented to me in exactly the same way as the minds of others. There are also lines of argumentation in contemporary analytical philosophy of mind that is close to and consistent with the phenomenological conception of pre-reflective self-awareness. Alvin Goldman provides an example: [Consider] the case of thinking about x or attending to x. In the process of thinking about x there is already an implicit awareness that one is thinking about x. There is no need for reflection here, for taking a step back from thinking about x in order to examine it…When we are thinking about x, the mind is focused on x, not on our thinking of x. Nevertheless, the process of thinking about x carries with it a non-reflective self-awareness (Goldman 1970, 96). A similar view has been defended by Owen Flanagan, who not only argues that consciousness involves self-consciousness in the weak sense that there is something it is like for the subject to have the experience, but also speaks of the low-level self-consciousness involved in experiencing my experiences as mine (Flanagan 1992, 194). As Flanagan quite correctly points out, this primary type of self-consciousness should not be confused with the

much stronger notion of self-consciousness that is in play when we are thinking about our own narrative self. The latter form of reflective self-consciousness presupposes both conceptual knowledge and narrative competence. It requires maturation and socialization, and the ability to access and issue reports about the states, traits, dispositions that make one the person one is. To claim that every kind of self-consciousness is conceptual is overly cognitive. Bermúdez (1998), to mention one further philosopher in the analytic tradition, argues that there are a variety of **nonconceptual forms of self-consciousness** that are "logically and ontogenetically more primitive than the higher forms of self-consciousness that are usually the focus of philosophical debate" (1998, 274; also see Poellner 2003). This growing consensus across philosophical studies supports the phenomenological view of pre-reflective self-consciousness. (References: Stanford Encyclopedia, Post Modersnist Theories, Internet Encyclopedia of Philosophy) Korzybski pointed out that there are two ways to slice easily through life; to believe everything or to doubt everything. Both ways save us from thinking. So, from time to time you have to give yourself a Virtual Reality Check. The hallmark of postmodern philosophy has been disbelief or skepticism of all "metanarratives," or translations of reality. Postmodernism has even turned its profound skepticism on such important humanist concepts as "objective truth" and reason.Yet, for a deconstructionist postmodern society, individually we are still riddled with superstition and gullibility, and open to manipulation through our belief systems as any politician, philosopher, clergy, or salesperson will attest. Further, most people are painfully naive when it comes to even the simplest scientific understanding. Most of us don't have a clue about the fundamental nature of physical reality or our own psychological nature, and our ability to be fooled by our senses and mind, or methods of social and technological persuasion. We may think we know, but the penetration is deeper than we dare imagine. Nature rejects the naiveté that seeks absolute truth. We are beginning to realize, individually and culturally, that "realities" are all human constructions. The task becomes one of "catching ourselves in the act" of creating our own "reality" from the flow of events. Human truth is always an engagement of mind with experience.. We don't need to fear the collapse of our

personalistic belief system (the "box" we live in), nor our belief in absolute truth. **A strong desire to engage in the "quest for uncertainty" complements our anxiety that perhaps there is no absolute, objective ground to reality**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t-T_8)F_8(R_8(t-T_8) - N_8(t)M_8((R_8(t) \qquad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8} + \frac{N_8 dG_8}{dR_8} - \frac{dB_8}{dR_8} > \frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8} + \frac{dM_8}{dR_8}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_8 = \frac{N_8 G_8}{2M_8} \frac{d(8nF_8)}{dR_8} (\geq 0) \quad (8a)$$

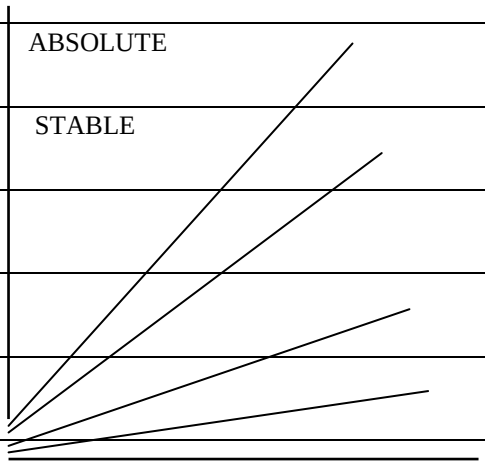
$$I_8 = -\frac{N_8 G_8}{2M_8} \frac{d(8nM_8}{dR_8} (\geq 0) \quad (8b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8 > L_8 - I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_8 > L_8 - I_8$, when $P_8^2/4 > I$ (8a)

$P_8 > L_8 - I_8 + \frac{(I_8 - P_8^2}{(L_8 - I)}$ when $P_8^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

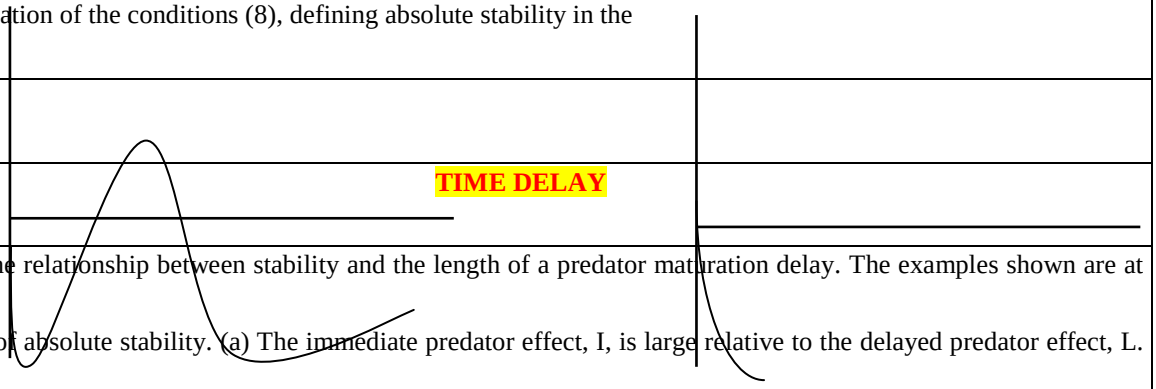



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this

rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior , we must note that the model could be applied to psychological systems too, albeit we are still not sure of the fact how the wave function collapses thereby due to perception.

Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility. He privatizes the sublime, having learned from the fate of his predecessors that the public can never be more than beautiful."(11) Notice is made here of what is indeed the telling feature. Levinas is praised, for example, for his alarm when "the social will [is] sought in an ideal of fusion . . . the subject losing himself in a collective representation, in a common ideal. . . . It is the collectivity which says 'us,' and which, turned toward the intelligible sun, toward the truth, experience, the other at his side and not face to face with him. The tellingly titled D'un ton apocalyptique contains passages that are particularly excited: These people situate themselves outside the ordinary, but they have in common this: they describe themselves as having an immediate and intuitive relation with mystery. And they want to attract, to seduce, and lead others to the mystery, through mystery. This agogic function of the leader of men, of the duce, of the Führer, of the leader, places him above the crowd that he manipulates with the aid of a small numbers of adepts joined together in a sect with a secret language, a clique or a small party with its

ritualized practices. The mystifiers pretend to have exclusive access to the privilege of a secret mystery. . . . The revelation or the unveiling of the secret is something that they jealously reserve for themselves. Jealousy is here a major characteristic. There is not narcissism and non-narcissism; there are narcissisms that are more or less comprehensive, generous, open, extended. What is called non-narcissism is in general but the economy of a much more welcoming, hospitable narcissism, one that is much more open to the experience of the other as other. I believe that without a movement of narcissistic reappropriation, the relation to the other would be absolutely destroyed; it would be destroyed in advance. The relation to the other--even if it remains asymmetrical, open, without reapporpriation--must trace a movement of reappropriation in the image of oneself for love to be possible, for example. Love is narcissistic. Beyond that, there are little narcissisms, there are big narcissisms. . . .(23) **Justice of the Pieces Deconstruction as a Social Psychology Douglas Collins** Sabine Wilke cogently argues for the rise of the aesthetic dimension in philosophical writing by showing how the metaphorical and poetic style of Heidegger, the principles of phrase configurations with Adorno, or the topographic writing of derrida replace through their very form something what traditionally would have been subject to discursive exposure The concern shifts from telling to showing, showing first and foremost that something remains untold. An exemplary book in this aesthetic drive is Avital Ronell's Telephone Book where the playful layout and explicit break with traditional design interrupt the sequential logic of the paragraphs thereby performing the very argument at stake: open up new circuits of meaning and signification. The engagement with form and publishing formats is only consequential since media theory as exposed by Walter J. Ong, Eric A. Havelock, Parry Miman or Marshall McLuhan besides many others made it clear that different communication technologies condition the mode of thought of a given culture The fractals of Benoit Mandelbrot or the butterfly shaped Lorenz attractor are the popular icons of this turn in the sciences. For communication studies and philosophy, Vilém Flusser was one of the pioneer thinkers reflecting the rise of what he called “techno-images.” The photograph is the paradigmatic case which condenses the whole crises of Western

thought and culture on its grainy surface. It subverts the linear logic of writing.6 Flusser's central diagnosis detects a shift from the grand narratives associated with the literate mind and Modernity towards creating and envisioning the world through image based practices. Writing himself in a succinct, visually inspired style, Flusser contents that the web of meaning, spun throughout the centuries by our textual practices, has become transparent. We discover the void between the letters through which the world escapes: "We are alienated from the world of text, precisely because we see through them as our own product."7 As a consequence we are confronted with a world of dissociated facts and information bits. Flusser hails this fragmented world as the advent of the zero dimensional point-universe that will provide the new raw material for our image based practices. Techno-images are nothing else than a conscious, inherently artistic drawing-together of point-elements as exemplified first and foremost by the silver-grains of the photograph but also by more recent pixel based techno-images. Similar to blueprints, designs or statistical curves, the weather map for example visualizes and contracts a tremendous amount of meteorological data from widely distributed stations into a single image of high- and low pressure

areas, temperatures, wind speeds, etc. Techno-images "compute" in this sense the universe of unimaginable, abstract information-bits into a concrete surface – the image. Echoing Paul Feyerabend, these image practices have to be understood as creative practices where the truth of scientific models has been replaced by its aesthetic quality **Aesthetics, Writing, Networked Computers Jörg Muller Dissertation Submitted to the Division of Media and Communications of The European Graduate School in Candidacy for the Degree of Doctor of Philosophy**

**Lothar Schäfer (1) and Sisir Roy (2) in** Quantum Reality, the Importance of Consciousness in the Universe, the Discovery of a Non-Empirical Realm of Physical Reality, and the Convergence with Ancient Traditions of Indian and Western Philosophy investigate metaphysical aspects that include 1) the discovery of a non-empirical part of physical reality in a realm of potentiality; 2) the emanation of the empirical world out of a realm of non-material forms; 3) the discovery that the nature of physical reality is that of an indivisible Wholeness – the One; and 4)

panpsychism: the possibility that the One is aware of its processes like a Cosmic Consciousness. The convergence of powerful traditions of seemingly disparate cultures is particularly important to point out in the present process of globalization, when a unifying view is needed to avoid controversy and conflict. To quote them in extenso : This is the nonlocality of the quantum world. Menas Kafatos and Robert Nadeau (1990) have drawn a remarkable conclusion from this phenomenon: If reality is nonlocal, the nature of the universe is that of an undivided wholeness. Because our consciousness has emerged from this wholeness and is part of it, it is possible to conclude that an element of consciousness is active in the universe: a Cosmic Consciousness. Pre-reflective self-consciousness is pre-reflective in the sense that (1) it is an awareness we have before we do any reflecting on our experience; (2) it is an implicit and first-order awareness rather than an explicit or higher-order form of self-consciousness. Indeed, an explicit reflective self-consciousness is possible only because there is a pre-reflective self-awareness that is an on-going and more primary self-consciousness. Although phenomenologists do not always agree on important questions about method, focus, or even whether there is an ego or self, they are in close to unanimous agreement about the idea that the experiential dimension always involves such an implicit pre-reflective self-awareness. In line with Edmund Husserl (1959, 189, 412), who maintains that consciousness always involves a self-appearance (Für-sich-selbst-erscheinens), and in agreement with Michel Henry (1963, 1965), who notes that experience is always self-manifesting, and with Maurice Merleau-Ponty who states that consciousness is always given to itself and that the word ‘consciousness’ has no meaning independently of this self-givenness (Merleau-Ponty 1945, 488), Jean-Paul Sartre writes that pre-reflective self-consciousness is not simply a quality added to the experience, an accessory; rather, it constitutes the very mode of being of the experience: This self-consciousness we ought to consider not as a new consciousness, but as the only mode of existence which is possible for a consciousness of something (Sartre 1943, 20 [1956, liv]). The notion of pre-reflective self-awareness is related to the idea that

experiences have a subjective 'feel' to them, a certain (phenomenal) quality of 'what it is like' or what it 'feels' like to have them.

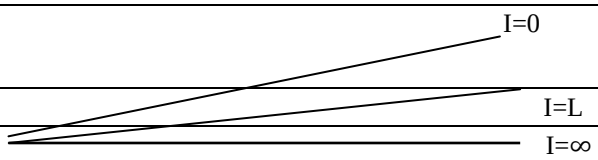
Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

TOTAL

ABSOLUTELY STABLE

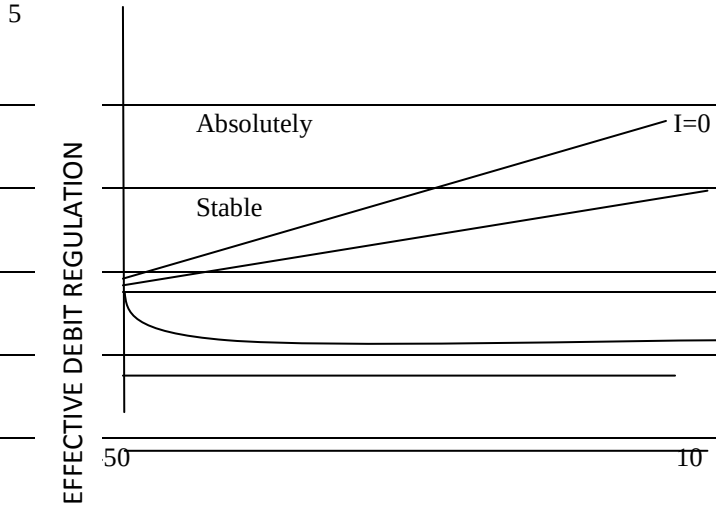



**DELAYED DEBIT EFFECT**

FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

## Utility Analysis And neural Network

Caroline Williams investigated affective density of the political and its effect on our understanding of political subjectivity. Taking up Spinoza's challenge to think about affect beyond corporeal embodiment, he argues that there is a modality of affectivity that cannot simply be inscribed within the borders of subjectivity. theorising affect as an impersonal force anchored in a relational ontology that gives due recognition to the circulation of affects, as well as to their ambivalent structure in creating sites of identification, and he utilises this ontology to reflect on the dynamic of the political and the shape of political subjectivity. Williams argues that Spinoza's philosophy (through ideas of conatus and imagination) offers the conceptual resources to reconfigure the composition of affective subjectivity as a transindividual social bond and as an unconscious dynamic of ethico-political existence. **Subjectivity (2010) 3, 245–262. doi:10.1057/sub.2010.15 Affective processes without a subject: Rethinking the relation between subjectivity and affect with Spinoza Caroline Williams** Phenomenological Time bears ample evidence and infallible observatory and apodictic evidence albeit with a sense of chagrined circumvention and defeated disgruntlment circularity of both psychological and phenomenological interpretations of art which are sometimes insightful, but most often work on the logic of cyclical time and unavoidable drives and representations played out as a ‘natural’ order: where ‘men can’t avoid being boys’ a circularity which valorizes a helplessness/truthfulness conundrum. One begins with a classical phenomenological fallacy with Duchamp’s Fountain, does the urinal keep returning or do we keep returning to the urinal? Put in this way, it is a male question, something that phenomenology conveniently conceals when positing the “I”, or indeed, the “we”. But we must be careful not to simply assume that the phenomenological first person is white, heterosexual, able-bodied and male, providing a foundation upon which knowledge should be based as a universal condition. Deleuze’s second critique of time is linked to Kant. This model of time extends not only to phenomenology but to the phenomenological interpretation

of art and is seen as a straight line. In the Critique of Pure Reason, Kant cuts up circular time and reassembles it as a series of sensory experiences which memory processes and which constitute subjectivity. Whereas the circular model of time (and interpretation) inscribes fatalism into the subject and the world, the phenomenological method is in danger of solipsism, where the subject constitutes the world. Much of embodied philosophy and enactivist theories of cognition go down this line with a plethora of art historical analyses by the likes of Rosalind Krauss, Michael Fried and others taking the lead from Merleau-Ponty. These approaches successfully challenge the **homunculus model of a Cartesian mind-body** dualism where there is a supervisor in the mind who processes experience and even consciousness and to whom all representations are directed. But phenomenology is mainly concerned with perceptual consciousness, the processing of colors, shapes, forms, motion in art and the world. The senses and sensorimotor processes are the fundamental pathways by which the interior is connected to the exterior in a chiasm of self-constitution in the world. The problem with this is that we explore the art through our sense perceptions, and what does art reveal to us?—our sense perceptions! Such solipsistic approaches have real problems relating complex and abstract conceptual production to bodily and sensorimotor contingencies, and they nearly always posit the subject as the site and object of discovery in the world. To be fair, phenomenology has moved on to rather more sophisticated models of intersubjectivity, taking into account the specificities of different kinds of bodies and ways of being beyond the white, male heterosexual as the common denominator, yet many naïve approaches to art perpetuate the notion that perceptual, rather than conceptual processes and evolutionary neural patterns are primary, fundamental ways to understand the art experience In fact, these aspects are only part of art production and reception, aspects which are often filtered out or made known to us in order for us to parody such trigger responses, as in much of Duchamp's work, which holds in contempt such retinal approaches to art. If we take the **subject of abstract expressionism** in the 1950s or later minimalist art, art history, teaching and interpretation still insists on putting the act of perceiving as fundamental to the constitution not only of the phenomenological first

person but also of art itself. The male act of usually standing up, which consists of directing a stream with the hand(s) into a container or on a flat surface, sand or snow in order to leave patterns and traces (the usual child's prank) underlies the logic of macho 1950s Abstract Expressionist culture, especially that of Jackson Pollock, so effectively parodied by Carolee Schneeman's Vagina Scroll, 1975. Here, painting or writing is a kind of a discharge which invests the **object and space with** the phenomenological "I", disciplined by existential struggle and reflected back by the traces, signatures and gestures left behind in the work of art, rediscoverable as an element in the making of other subjectivities.It is precisely this kind of solipsism which is blind to the art event as something well in excess of re-establishing habits of subjective self-identity through the production and interpretation of art. The attempt to reinforce a phenomenological identity upon the event by personalizing sense experience is unable to see art as something which remodels the phenomenological 'I' and neural plasticity itself. We learn new technologies and redraw our neural pathways through such new and spontaneous experiences which draw us into intersubjective worlds. Duchamp's Fountain and Deleuze's Repetition and Différence Grégory Minis sale (I have made some changes) The relation of different groups to deconstruction is settled by its relation to groups. An American enthusiast--one who has had a role in the development of our affection for Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^{\infty} \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$then given the constraints of $f(T_8)$, $(A1.5) can\ never\ be\ satisfied\ Co$nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8 \lambda_8 e^{-T_8 \lambda_8} + [(U_8 - V_8) M_8 + Y_8] - [((U_8 - V_8) M_8 -] e^{-T_8 \lambda_8} = 0$$

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$\boldsymbol{Z}_8^2 + M_8 Z_8 \sin(Z_8 T_8) - (A_8 M_8 + Y_8) + (A_8 M_8 + X_8)\cos(Z_8 T_8) = 0$$

$$(A_8 M_8) Z_8 - M_8 Z_8 \cos(Z_8 T_8) + (A_8 M_8 - X_8)\sin(Z_8 T_8) = 0$$

Where $A_8 = U_8 - V_8$Using (A2.3), substitute for $\boldsymbol{Z}_8^2$ and $Z_8$in (A2.2):

$W_8^2 (X_8^2 - 3 M_8 A_8 X_8 - M_8^2 (X_8 - Y_8) + 2 M_8^3 A_8 + 2 A_8^2 M_8^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_8$=1-cos $(Z_8 T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8 / 2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_8 \le 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2\ (\mathrm{X}_8 - Y_8)$$

$$Q_8(2, A_8) = \left(A_8 + 2\,\mathrm{M}_8\right)^2 (2A_8\,\mathrm{M}_8 - (\mathrm{X}_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2\left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8\left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8\left(2X_8 + Y_8\right)\right)$$

$$+2A_8\left(X_8 + Y_8\right)$$

$$=2\,M_8^2 W_8^2 (M_8 + \mathrm{A}_8) + M_8\, A_8^2 W_8 + \left(2A_8 M_8 - X_8 + Y_8\right)(A_8 W_8 + M_8 W_8^2$$

$$+ (2- W_8)((2\,M_8\,X_8 + M_8\,Y_8)\,W_8 + A_8\,(X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $\mathrm{W}_8$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in *A* will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range $0 \le W_8 \le 2$.Defining

$Q_8\left(W_8, A_8\right) = a_8 W_8^2 + c_8$ *where* $a_8\, b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \square - b_8 \square 0$$

And

$$b_8^2 \square 4a_8 c_8$$

$$Case(i): Y_8 = 0$$

$Let\ Y_8 = 0$ And try the substitution A= $X_8 = /2M_8 in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8 /$ $2M_8$) is positive for *W<2. Since* $Q_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\ X_8$ ,

$Let\ Y_8 = \ y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 + \ n_8 y_8)$ we have

$$W_8^2 y_8 \left[X_8 n_8 \frac{(1 + n_8 y_8)}{(2 + n_8 y_8)} + M_8^2 \left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8 X_8 \left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2}+2y_8\left(n_8-1\right)+\frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)}\right]=0.\ (A2.12)$$

Since $n_8$ and $A_8$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8=(2+n_8 y_8)$, *the condition*$(A2.88?)$*can be written as*

$[\frac{X_8(E_8-8)}{E_8^2}]^2[\frac{X_8^2}{M_8^2}-4y_8 E_8^2\frac{X_8}{M_8^2}+4y_8(n_8-2)E_8^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8/W_8^2)$ can be expressed as

$$2E_8[y_8E_8+(y_8E_8(8+y_8)(E_8-n_8))^{8/2}]>\frac{X_8}{M_8^2}>2E_8[y_8E_8-y_8E_8(8+y_8)(E_8-n_8))^{8/2}]$$

Thoughts on the Social and Political Implications of Correlationism .Lion Mirror Truth be told, as my thought has evolved the issue of correlationism had fallen off the radar for me. Somehow the debate had come to seem too "philosophical" to me, too "scholastic", too remote from what interests me: understanding why social assemblages are organized as they are, how power functions in social assemblages, and what we might do to address that power and change things. Somehow the question of whether or not we can get out of the correlation between thinking and being just came to seem remote from these sorts of issues. Somehow it seemed too epistemological. ASIDE: Numerous discussions over the years have led me to believe that the debate over correlationism is poorly understood (or maybe I just don't understand it). On countless occasions I've heard people say "of course we must relate to things in order to know them." Well yeah, of course! I don't think this is what the critic of correlationism is getting at. It seems to me that correlationism is something more robust than the theses that we must relate to something to know it.

Correlationism instead seems to require the theses that thought and being are indiscernible. Put more concretely, the correlationist is someone who argues that we either a) can never tell whether being is merely a construction of our thought (weak correlationism), or b) who argues that thought actually constructs being (strong correlationism). In other words, correlationism is another name for idealism. One can hold that we must relate to something in order to know it without being a correlationist. As an aside I should also add that I am a correlationist about some things. For example, I think money is something constructed by society and is therefore a strong correlationist when it comes to money. At any rate, for a long time I'd become rather indifferent to debates about correlationism and philosophies of access. I had learned the lessons of speculative realism– which I could have also learned, I think, from Deleuze and Guattari, the new materialist feminists, actor-network theorists such as Latour –and had moved on. However, occasionally you come across a tone of phrase that pitches something in a different light. In The Cut of the Real, Katerina Kolozova writes, …the political problem of contemporary philosophy identified by the 'new realists' is, in fact, the product of a more fundamental epistemic problem. In his book After Finitude, Quentin Meillassoux calls this problem 'correlationism' and identifies it as an essentially post-Kantian legacy, which continues to dominate and limit philosophy. As a matter of fact, correlationism lies at (eb) the heart of postmodern theory and consists in the premise that thought can only 'think itself,' that the real is inaccessible to knowledge and human subjectivity, and that there is nothing but discursive constructs that fully determine thinking and that are meth0dologically accounted for all the way down. (1 – 2) Thought thinking only itself: Thought only encountering itself. In the jargon of postmodern and poststructuralist lingo, this would be the thesis of infinite semiosis, where signs ("thoughts") only ever relate to other signs. Within this framework, discursivity comes to be the hegemonic framework defining all of being. At the level of politics and social theory more generally, if the correlationist thesis is true the

consequences are clear: all social phenomena are discursive and all solutions to social and political problems will be discursive. The sole sphere of the political will be the discursive and all questions of politics will be questions of speech-acts and interpretation. The problem here is not that many theorists recognize that the discursive and semiotic plays an important role in the social and the political. It does and I've repeated this tirelessly. The problem is with what happens when thought or the semiotic becomes a hegemon, an "all", foreclosing our ability to recognize other forms of power. What I've wanted to say is that not all power functions discursively. In my last post and elsewhere I spoke of some other forms of politics: Thermopolitics: The politics surrounding **energy in the form of calories and fuels such as gasoline** and coal, and how our life and our very bodies are structured by energy dependencies and by being trapped in particular distributive networks that render these forms of energy available. I'm being quite literal when I speak of energy, talking about the effects, for example, of the absence of food in certain educational environments on cognition, for example; and am generally hostile to metaphorical extensions of the concept of energy which I see as erasing the dimension of real materiality. Geopolitics: The role that features of natural and built geography such as mountain ranges, rivers, oceans, soil conditions, roads, housing design, etc., play in the form that social relations take and how they impact individual bodies. Chronopolitics: The way in which the structurations of time organize what is possible for us. For example, the structuration of the working day, how much we can say and comprehend at any given time, the impact of things like the invention of the clock, etc. **Oikopolitics: This would be the domain of political economy described so well by Marxists**. So five different types of politics: Semiopolitics (or what currently dominates critical theory), thermopolitics, geopolitics, chronopolitics, and oikopolitics. No doubt there are other sites of the political or political struggle that we could speak of, but this is a good start. Also, it should be obvious that these aren't exclusive domains, but are entangled in all sorts of important

ways. For example, something might take place at the level of semiopolitics (speech, law, rhetoric, norms, and communication) that has all sorts of impact at the level of thermopolitics. Congress might decide to cut programs that fund school meal programs. This, in turn, will have a thermodynamic impact on those students that go without the calories they need developmentally and cognitively to function in a particular way. There is an entanglement here of semiopolitical and thermopolitical domains. The young student here has been constrained both at the level of **semiotic phenomena and thermodynamic structures**. The point is that if true, semiotic intervention (speech-acts, protests, interpretations, deconstructions, etc) will not be an appropriate response to all political problems because social formations are not entirely structured by the semiotic. The child in that school does not suffer from a lack of the right signs, but from a lack of calories needed to run the engine of his thought and body. Certainly semiotic interventions might be needed to render that energy available, but it is the energy itself that is at issue and the absence of that energy that forms the spider web entangling him in his position. A correlationist perspective tends to erase this as even being a site of the political. **Larval Subjects January 28, 2014 Ontological Anarché Posted by larval subjects under Uncategorized January 27, 2014 Roden on Pluralism Over at enemy industry Roden has an excellent post up on the pluralism discussion. Check it out here. January 25, 2014 Different Senses of Pluralism and Ontology Posted by larval subjects under Uncategorized Responding to one of my comments over at his blog, Jeremy Trombley:** I have deleted restricted, reformulated some sentences .Kindly bear with me. Intention is to make the discursive enucleation consistent with the subject in question**.** Recourse to epistemology of the quantum mechanics generic and fervor of string theory is a sine qua non for the further comprehension of the theory and ontology .How would the world appear to us if its ontology was that of classical mechanics but every agent faced a restriction on how much they could come to know about the classical state? We show that in most respects

it would appear to us as quantum. The statistical theory of classical mechanics, which specifies how probability distributions over **phase space evolve** under Hamiltonian evolution and under measurements, is typically called Liouville mechanics, so the theory we explore here is Liouville mechanics with an epistemic restriction. The particular epistemic restriction we posit as our foundational postulate specifies two constraints. The first constraint is a classical analog of Heisenberg's uncertainty principle; the second-order moments of position and momentum defined by the phase-space distribution that characterizes an agent's knowledge are required to satisfy the same constraints as are satisfied by the moments of position and momentum observables for a quantum state. The second constraint is that the distribution should have maximal entropy for the given moments. Starting from this postulate, authors derive the allowed preparations, measurements, and transformations and demonstrate that they are isomorphic to those allowed in Gaussian quantum mechanics and generate the same experimental statistics. They argue that this reconstruction of Gaussian quantum mechanics constitutes additional evidence in favor of a research program wherein quantum states are interpreted as states of incomplete knowledge and that the phenomena that do not arise in Gaussian quantum mechanics provide the best clues for how one might reconstruct the full quantum theory. I: http://dx.doi.org/10.1103/PhysRevA.86.012103 **Reconstruction of Gaussian quantum mechanics from Liouville mechanics with an epistemic restriction Phys. Rev. A 86, 012103 – Published 10 July 2012 Stephen D. Bartlett, Terry Rudolph, and Robert W. Spekkens.** Deleuze in fact characterizes modes of existence, with their powers and capacities. The answer is this: Deleuze approaches modes of existence, ethically speaking, not in terms of their will, or their conscious decision making power (as in Kant), nor in terms of their interests (as in Marx, for example), but rather in terms of their drives. For Deleuze, conscious will and preconscious interest are both subsequent to our unconscious drives, and it is at the level of the drives that we have to aim our ethical analysis. Here, I would like to

focus on two sets of texts on the drives taken, not from Nietzsche and Spinoza, but rather from Nietzsche and Leibniz (Leibniz being one of the first philosophers in the history of philosophy to have developed a theory of the unconscious). The first set of texts comes from Nietzsche's great early book entitled Daybreak, published in July 1881. Nietzsche first approaches the question of the drives by giving us an everyday scenario: "Suppose we were in the market place one day," he writes, "and we noticed someone laughing at us as we went by: this event will signify this or that to us according to whether this or that drive happens at that moment to be at its height in us—and it will be a quite different event according to the kind of person we are. One person will absorb it like a drop of rain, another will shake it from him like an insect, another will try to pick a quarrel, another will examine his clothing to see if there is anything about it that might give rise to laughter, another will be led to reflect on the nature of laughter as such, another will be glad to have involuntarily augmented the amount of cheerfulness and sunshine in the world—and in each case, a drive has gratified itself, whether it be the drive to annoyance, or to combativeness or to reflection or to benevolence. This drive seized the event as its prey. Why precisely this one? Because, thirsty and hungry, it was lying in wait" (D 119). This is the source of Nietzsche's doctrine of perspectivism ("there are no facts, only interpretations"), but what is often overlooked is that, for Nietzsche, it is our drives that interpret the world, that are perspectival—and not our egos, not our conscious opinions. It is not so much that I have a different perspective on the world than you; it is rather that each of us has multiple perspectives on the world because of the multiplicity of our drives—drives that are often contradictory among themselves. "Within ourselves," Nietzsche writes, "we can be egoistic or altruistic, hard-hearted, magnanimous, just, lenient, insincere, can cause pain or give pleasure" (Parkes, pp. 291-292). We all contain such "a vast confusion of contradictory drives" (WP 259) that we are, as Nietzsche liked to say, multiplicities, and not unities. Moreover, these drives are in a constant struggle or combat with each other:

my drive to smoke and get my nicotine rush is in combat with (but also coexistent with) my drive to quit. This is where Nietzsche first developed his concept of the will to power—at the level of the drives. "Every drive is a kind of lust to rule," he writes, "each one has its perspective that it would like to compel all the other drives to accept as a norm" (WP 481). DELEUZE AND THE QUESTION OF DESIRE: TOWARD AN IMMANENT THEORY OF ETHICS Daniel W. Smith. **In on a Possible Physical Metatheory of Consciousness Miroljub Dugic et al** show that the modern quantum mechanics, and particularly the theory of decoherence, allows formulating a sort of a physical metatheory of consciousness. Particularly, the analysis of the necessary conditions for the occurrence of decoherence, along with the hypothesis that consciousness bears (more-or-less) well definable physical origin, leads to a wider physical picture naturally involving consciousness. This can be considered as a sort of a psycho-physical parallelism, but on very wide scales bearing some cosmological relevance. Open Systems & Information Dynamics Volume 9, Number 2, 153 (2002) arXiv: quant-ph/0212128. On a broader perspective and background, the theoretical computation of the universe is justified.IN physics and cosmology, digital physics is a collection of theoretical perspectives based on the premise that the universe is, at heart, describable by information, and is therefore computable. Therefore, according to this theory, the universe can be conceived of as either the output of a deterministic or probabilistic computer program, a vast, digital computation device, or mathematically isomorphic to such a device. Digital physics is grounded in one or more of the following hypotheses; listed in order of decreasing strength. The universe or reality: is essentially informational (although not every informational ontology needs to be digital); is essentially computable (the pancomputationalist position); can be described digitally; is in essence digital; is itself a computer (pancomputationalism); is the output of a simulated reality exercise. Following Jaynes and Weizsäcker, the physicist John Archibald Wheeler wrote the following: [...] it is not unreasonable to imagine that

information sits at the core of physics, just as it sits at the core of a computer. (John Archibald Wheeler 1998: 340) It from bit. Otherwise put, every 'it'—every particle, every field of force, even the space-time continuum itself—derives its function, its meaning, its very existence entirely—even if in some contexts indirectly—from the apparatus-elicited answers to yes-or-no questions, binary choices, bits. 'It from bit' symbolizes the idea that every item of the physical world has at bottom—a very deep bottom, in most instances—an immaterial source and explanation; that which we call reality arises in the last analysis from the posing of yes–no questions and the registering of equipment-evoked responses; in short, that all things physical are information-theoretic in origin and that this is a participatory universe. (John Archibald Wheeler 1990: 5) David Chalmers of the Australian National University summarised Wheeler's views as follows: Wheeler (1990) has suggested that information is fundamental to the physics of the universe. According to this 'it from bit' doctrine, the laws of physics can be cast in terms of information, postulating different states that give rise to different effects without actually saying what those states are. It is only their position in an information space that counts. If so, then information is a natural candidate to also play a role in a fundamental theory of consciousness. We are led to a conception of the world on which information is truly fundamental, and on which it has two basic aspects, corresponding to the physical and the phenomenal features of the world. (Wikipedia). **Participatory consciousness forms the bastion, pillar, post, stylobate and fulcrum of the redressal measures and corrective steps taken for circumvention of many problems, identified and unidentified, spread widest commonalty. Anne Baring's sententious and pithy article will be good to have a timely outlook.** I would like to open my contribution to this meeting with the words of Ken Wilber that I quoted in the first Millennium Symposium: 'The secret impulse of life is towards greater consciousness. Maybe the evolutionary sequence is from matter to body, to mind, to soul, to spirit - each transcending and including, each with a greater depth and greater

consciousness and wider embrace' Ken Wilber, and with these from Rupert Sheldrake's contribution to the same meeting: "It is possible to make quite a good case that the sun could be conscious and if the sun, then why not the stars? Why not the galaxies? Why should there not be galactic minds? And if they communicate with each other then the distances are far too great for light to be the medium of transmission. We are into inter-galactic telepathy." After many decades of study, I know of only one diagram in all religious traditions that can adequately explain these statements; that can give us both a blue-print of the many levels of consciousness and a route-map to help us reach beyond our present level of understanding. I would like to introduce a diagram of the four worlds of Kabbalah to illustrate the theme that spirit has brought these four worlds into being and that we, living in the fourth and manifest world know nothing of the existence of the other three invisible ones. It is the clearest diagram I know of that can explain to us the multi-levelled structure of consciousness and our place in the great chain of being. We are indeed, as Wilbur suggests, evolving from matter to body, to mind, to soul, to spirit but where have we come from? It is possible that, as participants in the emanation of spirit into manifest life, we have also come from spirit to soul, to mind, to body, to matter. Everything, from the immense galactic energies to the minute forms of matter participates in one life, is generated from one creative spirit. There is nothing outside this life, this spirit. This diagram suggests that truth is not an article of faith but the experience of participating in the undiscovered reality of what has brought us into being. We cannot know truth until we have entered into communion with that reality. Continuing in similar vein **Michel Bitbol dilates upon the dialectic deliberation and polemical argumentation between** hermeneutists and eliminativists. When he formulated the program of Neurophenomenology, Francisco Varela suggested a balanced methodological dissolution of the “hard problem” of consciousness. He shows that his dissolution is a paradigm which imposes itself onto seemingly opposite views, including materialist approaches. I also point out that

Varela's revolutionary epistemological ideas are gaining wider acceptance as a side effect of a recent controversy between hermeneutists and eliminativists. Finally, he emphasizes a structural parallel between the science of consciousness and the distinctive features of quantum mechanics. This parallel, together with the former convergences, point towards the common origin of the main puzzles of both quantum mechanics and the philosophy of mind: neglect of the constitutive blindspot of objective knowledge. Phenomenology and the Cognitive Sciences Phenomenology and the Cognitive Sciences **Phenomenology and the Cognitive Sciences 2002, Volume 1, Issue 2, pp 181-224 Science as if situation mattered Michel Bitbol**

**ABSOLUTE STABILITY**

Hence if $X_8 / M_8^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_8$ ) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8 - 8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2 - 2y_8\frac{(E_8 + n_8^2 y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right) - 4y_8(E_8 - n_8)$$

And has the solution

$$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8 - 8)}[y_8(E_8 + n_8^2 y_8) + [y_8^2 (E_8 + n_8^2 y_8)^2 + 4 y_8 (E_8 - 8)(E_8 - n_8)]^{8/2}]$$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2E_8[y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSOILIDATION**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + (M_8 + U_8)\lambda_8 - (M_8 + V_8)\lambda_8 e^{2T_8\lambda_8} + (M_8 U_8 + Y_8)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8\, e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + (M_8 + V_8) Z_8 \sin(Z_8 T_8) - (U_8 M_8 + Y_8) + (M_8(U_8 + V_8) - X_8)$$

$$\mathrm{X}\cos(Z_8 T_8) - 2M_8 V_8 \cos^2(Z_8 T_8) + M_8 V_8 = 0$$

$$(M_8 + U_8) Z_8 - (M_8 + V_8) Z_8 \cos(Z_8 T_8) + (M_8(U_8 + V_8) - X)$$

$$\mathrm{X}\sin((Z_8 T_8) - M_8 V_8 \sin(Z_8 T_8) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 =$

$1 - \cos Z_8\, T_8$and $A_8 = U_8 - V_8$we obtain

$$4M_8^2 V_8^2 W_8^4 + 2M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8(M_8 + V_8) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2[M_8 - V_8)^2(2A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8(M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8(A_8^2 + 4(X_8 + Y_8))]$$

$$+ W_8 \left[A_8^2 + 2X_8\right)\left(2A_8 M_8 - X_8\right) + 2A_8 Y_8 (M_8 + V_8) - M_8 A_8^3] + A_8^2 \left(X_8 M_8 - Y_8\right) = 0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8 > 0$$

$$and$$

$$C_8 > (X_8 - Y_8)/2M_8$$

$recalling\ that\ A_8 = U_8$ and defining $C_8 = U_8 + V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8(0, U_8 V_8 = A_8^2 (X_8 + Y_8)$$

$$Q_8\left(2, U_8 V_8\right) = \left(C_8 + 2M_8^2\right)^2 (2C_8 M_8 + Y_8 - X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$ the existence of a D-position at this point is determined by letting cos $(Z_8 T_8) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all $T$, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**ABSOLUTE STABILITY IN PREDATOR- PREY MODELS**

**APPLIED TO ACCETUATION ATTRITION MODELS**

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability

criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The Percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending

upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t-T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\big(R(t)\big)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The

regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the

TIME DELAY

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a dominant eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

DELAY DUE TO CONGREGATIONAL AND CONFEDERATIONAL DEBIT CREDIT TYEPES

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t-T_8)F_8(R_8(t-T_8) - N_8(t)M_8((R_8(t) \quad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8} + \frac{N_8 dG_8}{dR_8} - \frac{dB_8}{dR_8} > \frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8} + \frac{dM_8}{dR_8}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_8 = \frac{N_8 G_8}{2M_8}\frac{d(8nF_8)}{dR_8} (\geq 0) \quad (8a)$$

$$I_8 = -\frac{N_8 G_8}{2M_8}\frac{d(8nM_8}{dR_8} (\geq 0) \quad (8b)$$

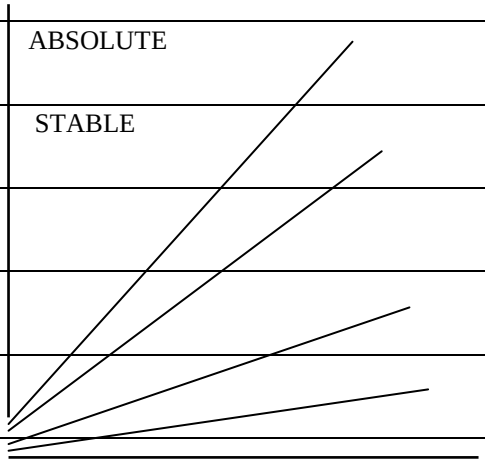
 

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8 > L_8 - I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_8 > L_8 - I_8, \quad \text{when} \quad P_8^2/4 > I \tag{8a}$$

$$P_8 > L_8 - I_8 + \frac{(I_8 - P_8^2}{(L_8 - I)} \quad \text{when} \quad P_8^2/4 < I. \tag{8b}$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit

(like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

See the diagram Nunney gives for predator and prey analysis. We can simulate the same for the financial system in the foregoing too.

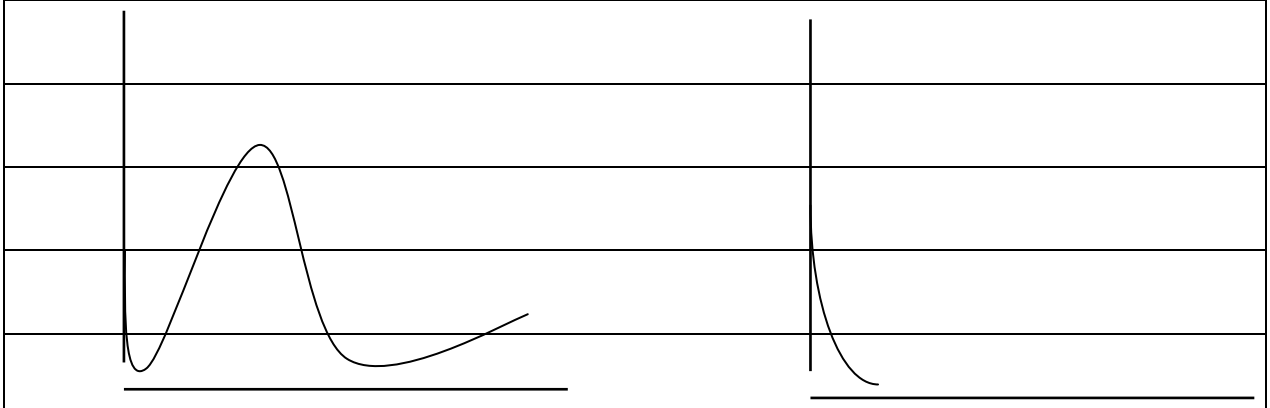
Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the

debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

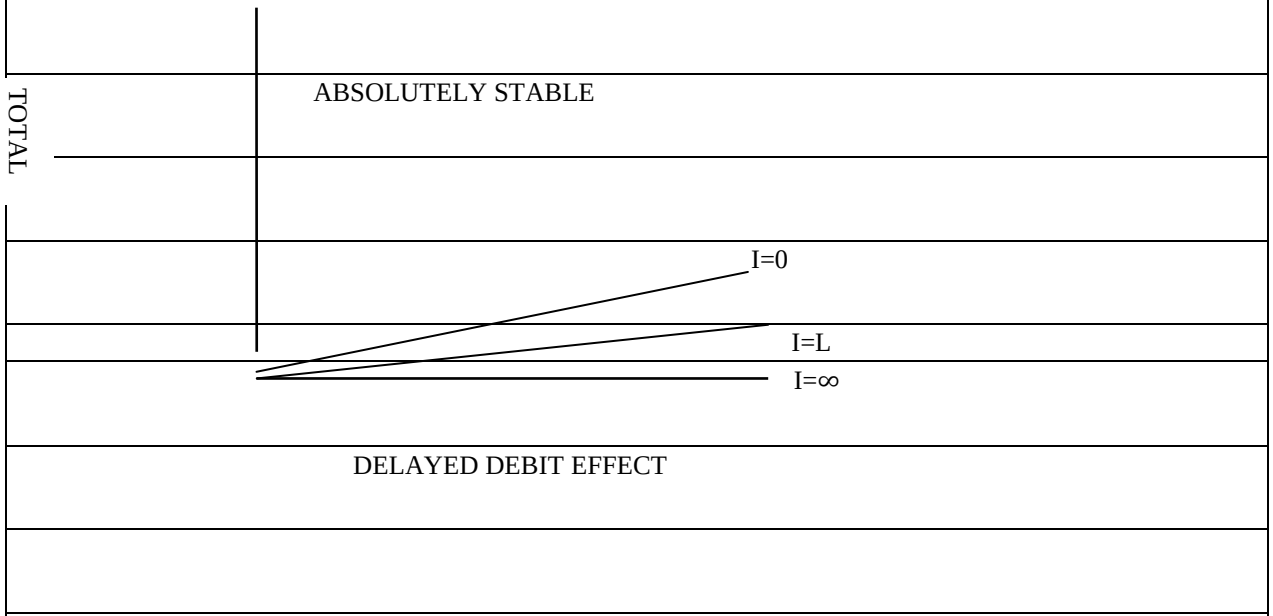
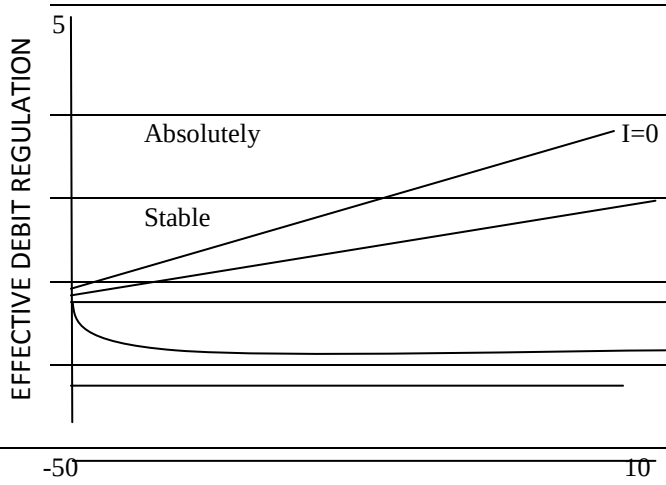



FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison?

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are

defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^{\infty} \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$ then given the constraints of $f(T_8), (A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8 \lambda_8 e^{-T_8 \lambda_8} + [(U_8 - V_8) M_8 + Y_8] - [((U_8 -$$

$$V_8 M_8 -] e^{-T_8 \lambda_8} = 0$$

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$\boldsymbol{Z}_8^2 + M_8 Z_8 \sin(Z_8 T_8) - (A_8 M_8 + Y_8) + (A_8 M_8 + X_8) \cos(Z_8 T_8) = 0$$

$$(A_8 M_8) Z_8 - M_8 Z_8 \cos(Z_8 T_8) + (A_8 M_8 - X_8) \sin(Z_8 T_8) = 0$$

Where $A_8 = U_8 - V_8$ Using (A2.3), substitute for $\boldsymbol{Z}_8^2$ and $Z_8$ in (A2.2):

$W_8^2 (X_8^2 - 3 M_8 A_8 X_8 - M_8^2 (X_8 - Y_8) + 2 M_8^3 A_8 + 2 A_8^2 M_8^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_8 = 1 - \cos(Z_8 T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8 / 2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_8 \le 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2\ (\mathrm{X}_8 - Y_8)$$

$$Q_8(2, A_8) = \left(A_8 + 2\,\mathrm{M}_8\right)^2 (2A_8\,\mathrm{M}_8 - (\mathrm{X}_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2\left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8\left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8\left(2X_8 + Y_8\right)\right)$$

$$+2A_8\left(X_8 + Y_8\right)$$

$$=2\,M_8^2 W_8^2 (M_8 + \mathrm{A}_8) + M_8\, A_8^2 W_8 + \left(2A_8 M_8 - X_8 + Y_8\right)(A_8 W_8 + M_8 W_8^2$$

$$+ (2- W_8)((2\,M_8\,X_8 + M_8\,Y_8)\,W_8 + A_8\,(X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $\mathrm{W}_8$touches the $W$ –axis only at its minimum value in which case (A2.9) shows that further increases in *A* will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range $0 \le W_8 \le 2$.Defining

$Q_8\left(W_8, A_8\right) = a_8 W_8^2 + c_8$ *where* $a_8\, b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \square b_8 \square 0$$

And

$$b_8^2 \square 4a_8c_8$$

$$Case(i): Y_8 = 0$$

$Let\ Y_8 = 0$ And try the substitution A= $X_8 = /2M_8 in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8 /$ $2M_8$) is positive for *W<2. Since* $Q_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\ X_8$ ,

$Let\ Y_8 = \ y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 + \ n_8 y_8)$ we have

$$W_8^2 y_8\left[X_8 n_8 \frac{(1 + n_8 y_8)}{(2 + n_8 y_8)} + M_8^2\left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8 X_8 \left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2}+2y_8\left(n_8-1\right)+\frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)}=0.(A2.12)\right.$$

Since $n_8$ and $A_8$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8=(2+n_8 y_8), the\ condition (A2.88?) can\ be\ written\ as$

$[\frac{X_8(E_8-8)}{E_8^2}]^2[\frac{X_8^2}{M_8^2}-4y_8 E_8^2\frac{X_8}{M_8^2}+4y_8(n_8-2)E_8^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8/W_8^2)$ can be expressed as

$2E_8[y_8E_8+(y_8 E_8(8+y_8)(E_8-n_8))^{8/2}]>\frac{X_8}{M_8^2}>2E_8[y_8E_8-y_8E_8(8+y_8)(E_8-n_8))^{8/2}]$

**ABSOLUTE STABILITY**

Hence if $X_8/M_8^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0 $W_8$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8-8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2-2y_8\frac{(E_8+n_8^2y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right)-4y_8(E_8-n_8)$$ 0

And has the solution

$$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8-8)}[y_8(E_8 + n_8^2\, y_8) + [y_8^2\,(E_8 + n_8^2\, y_8)^2 + 4\, y_8 (E_8 - 8)(E_8 - n_8)]^{8/2}]$$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2\, E_8 [y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + (M_8 + U_8)\lambda_8 - (M_8 + V_8)\lambda_8 e^{2T_8\lambda_8} + (M_8 U_8 + Y_8)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8\, e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + (M_8 + V_8) Z_8 \sin(Z_8 T_8) - (U_8 M_8 + Y_8) + (M_8(U_8 + V_8) - X_8)$$

$$\mathrm{X}\cos(Z_8 T_8) - 2M_8 V_8 \cos^2(Z_8 T_8) + M_8 V_8 = 0$$

$$(M_8 + U_8) Z_8 - (M_8 + V_8) Z_8 \cos(Z_8 T_8) + (M_8(U_8 + V_8) - X)$$

$$\mathrm{X}\sin((Z_8 T_8) - M_8 V_8 \sin(Z_8 T_8) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 = 1 - \cos Z_8 T_8$and $A_8 = U_8 - V_8$we obtain

$$4M_8^2 V_8^2 W_8^4 + 2M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8(M_8 + V_8) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2[M_8 - V_8)^2(2A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8(M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8(A_8^2 + 4(X_8 + Y_8))]$$

$$+ W_8[A_8^2 + 2X_8)(2A_8 M_8 - X_8) + 2A_8 Y_8(M_8 + V_8) - M_8 A_8^3] + A_8^2(X_8 M_8 - Y_8) = 0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8 > 0$$

$$and$$

$$C_8 > (X_8 - Y_8)/2M_8$$

$recalling\ that\ A_8 = U_8$ and defining $C_8 = U_8 + V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8(0, U_8 V_8 = A_8^2 (X_8 + Y_8)$$

$$Q_8\left(2, U_8 V_8\right) = \left(C_8 + 2M_8^2\right)^2 (2C_8 M_8 + Y_8 - X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$the existence of a D-position at this point is determined by letting cos $(Z_8 T_8) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all *T*, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

## ABSOLUTE STABILITY IN PREDATOR- PREY MODELS

## APPLIED TO ACCENTUATION ATTRITION MODELS

We apply model to each and every module and a concatenated model was to be presented which we procrastinate due to spatial restrictions.

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - D(R) - NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

The question is political-discursive. It's what Stengers and Latour are getting at with their concepts of diplomacy and cosmopolitics. They grant, first, that all entities exist and, second, that to say that someone's cherished idol (or whatever disputed entity they hold dear) is non-existent is a **'declaration of war'** – 'this means war,' as Stengers often says. They thus shunt onto-political discourse off of the terrain of knowledge/belief in the sense of existence/non-existence. Their basic claim seems to be that 'respect for otherness,' i.e. political pluralism, can only come from granting the entities that others hold dear an ontology, even if you don't 'believe' in them. You

are thus permitted to say 'I do not follow that god, he has no hold over me' but you are not permitted to say 'your god is an inane, infantile, non-existent fantasy, grow up.' And it's not just a question of politeness (although there's that too). The point is to grant others' idols and deities an existence – one needn't agree over what that existence entails, over what capacities that entity has or what obligations it impresses upon you as someone in its partial presence but to deny it existence entirely is to 'declare war' – to deny the possibility of civil discourse, of pluralistic co-existence. I believe that it was Richard Rorty who once quipped something like claims to reconcile realism and idealism always seem to end with a triumph of idealism. We don't, in fact, get realism through such approaches, but rather just get a pervasive anti-realism. I think this is also the problem with the "non-controversial pluralism" advocated by Stengers and Latour that Phillip defends here. Such pluralism is not realism but is, in fact, a thoroughgoing social constructivism. I think this is the central problem with Latour's argument in Irreductions (these days I regret having ever defended it). In rejecting both Enlightenment critique and what he calls "reduction" he wants to say something like "The Pentacostal really is filled with (e&eb) the Holy Spirit", that for the 19th working scientist heat really is a fluid and phlogiston really is what allows things to burn, and that for the Greek lightning really is an expression of Zeus's anger. Latour tells us that we aren't to reduce or explain away the entities posited by another group's "ontology" but are to develop explanations from within that ontology. It might not sound particularly sexy– and it certainly doesn't tell us what is worth thinking –but I can't help but believe that philosophy is the critical and reflective investigation of basic concepts that guide our investigation of the world about us, how we ought to live our lives, and what form of governance might be best. Compare two figures. A scientist might ask, what causes depression? We can very well imagine a philosopher turning around and asking the scientist, what is causality? The scientist presupposes a concept of causality in her investigations. She uses this concept in her inquiry. Now she might have a sophisticated concept of causality or she might never have thought much about causality at all, using it in the sort of colloquial and unreflective way that Plato decried when, for example, people

like Euthyphro talked about piety. A whole cascade of questions arise when we raise a question like "what is causality?" We can ask whether or not causality exists at all. We can ask how we distinguish between correlation and two events that merely accompany one another from genuine causation. This, for example, was Hume's question. But perhaps most importantly we can ask whether there is only one form of causality or many forms of causality. Is there only one-to-one causation; one cause and one effect? Is there many-to-one causation; or many events conspiring to produce an effect? Is there one-to-many causation; or one event producing a variety of different effects? We can even ask whether causality necessarily moves from past to present or whether there aren't forms of causality that move from future to past! **January 23, 2014 Pluralism and Realism:** multiple-worlds Over at Struggles Forever, Jeremy Trombley has an interesting post up on "the ontological turn" in anthropology or ethnography. I've been meaning to have a discussion with him about this as I think it's an issue many of us are struggling with. For example, the core project of The Democracy of Objects– a project which I think many have missed –is to somehow reconcile some version of social constructivism with a realist ontology capable of making room for ecology (which requires realist and materialist positions as there's a fact of the matter where global warming is concerned) as well as the role played by objective agencies in social assemblages such as technologies, infrastructure, features of geography, local climates, the growth cycles of plants and animals, waste, etc. Maybe we can try to organize some cross-blog event to discuss these issues. I certainly think they're close to the heart of Jeremy, Michael of Archive Fire, Arran James, and a host of others. As an aside, I'm beginning to realize how the different sites of the political I've been outlining– semiopolitics, thermopolitics, oikopolitics (political economy), geopolitics, eropolitics (the politics of sex and desire), biopolitics, and chronopolitics (and I'm sure there are other political sites!) –are drawing me away from traditional Marxism. Assuming that classical Marxism holds that economics or the conditions and relations of production are determinative of all other sites of the political, the various sites of the political that I've been outlining would lead to the conclusion that there is not one determinative

base of the political. This would not require committing Marx to flames, but rather of recognizing the phenomenon of overdetermination, or of a variety of different entangled sites of the political. But I digress. First, I find myself wondering what the ontological turn means in ethnography. Is it 1) the investigation of the different ontologies held/proposed by different cultures? E.g., the Aztecs believed that reality was structured in this way, while the Greeks in that way, and the ancient Chinese this way, etc? Or 2) is it an investigation of how real entities– independent of cultural beliefs –influence cultural formations? Or is it a combination of both? A position that I would favor. January 20, 2014 Model is applied to structured reality

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

**The evaporation spectrum of black holes from a local quantum gravity perspective Aurélien Barrau (Submitted on 21 May 2016**)

(11) Authors revisit the hypothesis of (e) a possible line structure in the Hawking evaporation spectrum of black holes, due to (e) non-perturbative quantum gravity effects, even arbitrarily far away from (e) the Planck mass.

(12) Authors show that this naive prediction might in fact hold in (eb)the specific context of loop quantum gravity, with (e&eb) a small departure from (e) the ideal case for (e) some low-spin transitions.

They also show that the effect is (=) neither washed out by the dynamics of (e) the process, nor by (e) existence of a mass spectrum up to (e) a given width, nor by the secondary component induced by (e) the decay of neutral pions emitted during (e&eb) the time-integrated evaporation. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Cite as: arXiv:1605.06613 [gr-qc] (or arXiv:1605.06613v1 [gr-qc] for this version

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t - T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\left(R(t)\right)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

It's very easy to treat this as an abstract, academic debate: "Are you a realist or are you an anti-realist?", as if it were just a matter of what happens to be true. But it seems to me that this debate has, in reality, always been about politics. As I outlined in my last post, we have perpetually seen how appeals to the real and natural have been used in the name of oppressive power, inscribing both the exploitation of nature and the oppression of various people in the very fabric of being itself. Theistic theology and realist ontology have perpetually been used in the name of

what Deleuze called "State Philosophies" or philosophies that ontologize contingent orders of power and privilege (e.g., "the great chain of being" used to justify patriarchy, monarchy, serfdom, poverty, etc, and appeals to nature used to justify poverty and racial inequality (The Bell Curve), patriarchy (evolutionary psychology), heteronormativity, capitalism, etc). Because arrangements of power and inequality are always contingent in the sense that there's no marked difference in the capacities of peoples, power always looks for a transcendent supplement that would provide justification through ontological necessity. Antirealism– from the Greek atomists to present –became the radical and emancipatory gesture because it revealed the lie behind all of these forms of social organization or their inherit contingency or arbitrariness. Realism, by contrast, has all too often functioned as an apologetics for arbitrary power and social organizations. Here it's worth recalling what Foucault said about science: "…Even before we know to what extent something like Marxism or psychoanalysis is analogous to a scientific practice in its day-to-day operations, in its rules of construction, in the concepts it uses, we should be asking the question, asking ourselves about the aspiration to power that is inherent in the claim to being a science. The question or questions that have to be asked are: "What types of knowledge are you trying to disqualify when you say that you are a science? What speaking subject, what discursive subject, what subject of experience and knowledge are you trying to minorize when you say 'I speak this discourse, I am speaking a scientific discourse, and I am a scientist.' What theoretico-political vanguard are you trying to put on the throne in order to detach it from the massive, circulating, and discontinuous forms that knowledge can take?" (Society Must Be Defended, 10). All of these questions hold equally for claims to something being real. What is one trying to minorize when claiming something is real? What becomes privileged? What is excluded? It is these questions that have been at the heart of the realism debates, for as Spencer-Brown taught us, every distinction has a marked and unmarked space, draws attention to something to be included and pushes something into the unconscious or the domain of the invisible, hidden, or veiled. This is above all the case with evocations of the real. However, as I've tried to show antirealism

leads to its own problems. First, so long as we exclude real beings from our ontological inventory, we are unable to fully understand how power functions ("such and such a set of cultural formations have been thoroughly debunked, yet people still live as if they believed them"). Not only do we not fully understand the sources of the problems due to too much focus on the discursive and semiotic, but we deny ourselves valuable sites of political intervention at the level of infrastructure. Second, we are prevented from addressing things such as cultural racism, such as that found in Heidegger's privileging of the West and the Greeks and Germans in particular. We do a good job addressing biological and theological racism, heteronormativity, and sexism by showing how it is a cultural construction, but when faced with racism such as Heidegger's where he argues that there's something "unique" about the Greek event, it's language, and about the German language, or Badiou/Zizek's racism with respect to the "Pauline Event", we really have no response. Here someone like Jared Diamond or Fernand Braud el is needed to explain global-geographical inequalities. Third, the tools of the cultural turn really do not provide us with the means of thinking the ecological as a site of the political. For this reason, I've tried to formulate a third way– which might be called "constructivist realism" or "constructivist naturalism" –that retains the insights of the cultural turn, while also allowing a robust place for the material. I don't claim to be original in this. I think that many such as Manuel DeLanda, Deleuze and Guattari, Stacy Alaimo, Karen Barad, Jane Bennett, etc., are up to something similar. It's a vast project that requires the work of a multitude of voices, especially given the way in which the culturalism pervades humanities. **The Relational and the Non-Relational: Notes towards an Immanent and Pluralist Theory of Meaning Posted by larval subjects under uncategorized k-bigpic I have deleted or reformulated some sentences for which I beg pardon attributable to spatial constraints** What is correlation, what is causation? What is coordination? What is coincidence? These thoughts have bothered humans since long. I have been able to excavate some reliable material from Wikipedia which I would like to share. "All views are partial and contingent – that's the lesson of pluralism, and you say as much yourself when you say that we must be attentive to ways our own

knowledge might contain superstitions. But it's not discourse that determines whether one view works and another doesn't – it's a confrontation with **non-discursive (non-human)** agencies." I wonder if part of the issue here is that we understand pluralism differently. For me this doesn't sound like pluralism at all. Rather, it just sounds like our **epistemological condition**. We [hopefully] come to understand those portions of existence we question or investigate. The reason we come to investigate them is largely contingent. Finally, the accounts of these features of existence we give can be mistaken. All of this is perfectly consistent with a monism. I'm not sure why what he outlines above is the lesson of pluralism. It seems to me that every realist knows this. It seems to me that ontological pluralism is something quite different. Ontological pluralism is the thesis that there are many different worlds inhabited by many different entities. Thus, for example, you would have one world; say that of Lucretius, that's only inhabited by atoms and their combinations. You would have another world; say the world of the Mongolian shaman, that's inhabited by spirits that do all sorts of things. The ontological pluralist is saying that these spirits are, that they exist, that they're real. This is something quite different than merely saying that there are partial and contingent points of view on the world. It's this saying that these entities are rather than that some person or groups of people believe that they are that is the nub of the issue. Now, I think part of the issue here is that there's an ambiguity in the term ontology. Ontology can be one of two things. On the one hand, ontology is a group or persons set of beliefs as to what is. Here it's trivially true that there are pluralities of ontologies and the realist readily recognizes this. This is the whole reason there are debates over ontology. Mongolian shamans have their ontology, Europeans theirs, Christian fundamentalists theirs, materialists theirs, etc. When striving to understand and communicate with others it's vital to understand these ontologies because, as rhetoricians like Burke point out, our **beliefs about what is** are among the things that motivate our action. Despite having never seen bacteria I was my hands because I believe there are bacteria and viruses on door handles and whatnot and don't want to get sick. My belief about a particular thing existing is what motivates my action. And who knows, perhaps this belief is as

superstitious as the belief that the crops failed because God was displeased with my community. On the other hand, ontology is a theory **about what is.** It is making a claim that something exists. This is where the rubber hits the road. The ontological pluralist seems committed to the thesis that every group set of beliefs about what exists is sufficient for granting the existence of those entities. This is what I find objectionable in Latour. Obviously I'm not bothered by Latour's suggestion that we should take into account the role that nonhumans like speed bumps, rivers, microbes, etc., play in the form that social assemblages take. But this is not all that Latour claims. This is precisely where the philosopher might balk. Unlike the ethnographer, the philosopher is not interested in what people believe exists, but rather philosophers– at least of the realist variant –are trying to figure out what is. In other words, the realist philosopher begins with the premise that not all of these beliefs about what is are true. So for the philosopher, recognizing that Mongolian shaman's believe in the existence of shamans would only be the first step. The next step would consist in determining whether there's good reason for thinking such entities really do exist, i.e., whether there's good reason for believing these entities have mind and culture independent reality. Lest readers think that I'm just picking on the supernatural here, we can ask similar questions about strings, subatomic particles, galaxies, etc., etc., etc. I guess the question really comes down to what exactly we mean by ontological pluralism. When we talk about ontological pluralism are we defending the thesis that people have different theories of being? If so, then ontological pluralism is trivially true. If this is what is meant, then I certainly share Jeremy's view that it's valuable to understand the different world's people believe in. Certainly when I was practicing as an analyst I didn't get in ontological debates with my patients and it was necessary to understand their theory of being or their ontology to properly attend to them. Or, when we talk about ontological pluralism are we defending the thesis that all these ontological theories are true and refer to really existing entities? That's quite a different claim and is not one I would defend or endorse. Now someone might object that "in both The Democracy of Objects and Onto-Cartography you defend the thesis that there are multiple worlds." This is true. Because I hold that not everything

is related I'm led to the conclusion that there are diverse worlds. However, I also hold that however many worlds there might be, these worlds are nonetheless composed solely of material entities. Within the framework I propose I wouldn't suggest that there's one world where there are spirits and another world where there are souls and yet another where there are only material entities. My view is that there aren't spirits or souls in any of these worlds. One of the things I keep hearing in these discussions is that somehow the realist adopts a view from nowhere. I honestly don't understand this criticism. Investigation always occurs somewhere and requires all sorts of mediations involving technologies, experiments, etc. It's that labor of gathering evidence, conducting experiments, using technologies to observe the world, etc., that gradually gives us a body of data that allows us to say there's good reason to believe that such and such a thing exists and has these powers. Another charge seems to be that the realist refuses to recognize that their claims about the world are fallible. I find this charge particularly strange because it's precisely because the realist recognizes the difference between our theories of the world or what we say about the world and the world itself that fallibility is built into the core of his position. Realism doesn't mean one holds they have special access to the world, that they know all truths, or that they have the truth in hand, only that there are truths to be known and that we can be mistaken about things. Circling **Squares Philip** has a post responding to my quandaries about how to mesh realism and pluralism. He writes: Ontologically and metaphysically the idea of realist pluralism is no longer an issue. There are (appropriately) numerous variants but the basic idea that reality is itself pluralistic is well established. Captioned expatiation is about Everett worlds and string theory

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from

being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

The von Neumann formulation leads us to conclude that the lower realms of quantum reality form a virtually limitless array of potentials that are inherently incapable of realization without the observation of a knowing entity, which we identify with consciousness (Kafatos and Nadeau, 1990). We thus have a correspondence between quantum theory, the hierarchical structure implied by the Holographic Principle, and the higher orders of experience,

leading to the full expression of the Conscious Universe. David Bohm's theory of the implicate and explicate orders involves a holographic Principle that is fully consistent with the Holographic Principle discussed here (Bohm 1980; 1986; Bohm and Hiley, 1993). According to Bohm, there is an implicate order that represents the universal, holographic subtext of reality, and which unfolds in every moment to produce the explicate order that we all observe. Thomas Germinario (2004) has equated the implicate order with (eb) the unconscious process, and the explicate order with (eb) conscious process. He emphasizes the nature of dreams within (eb) the implicate order, and the importance of the dream work for maintaining a healthy mind through (e&eb) integration of the implicate order and our daily lives. Allan Combs and Mark Holland (1990) connected the implicate order or holomovement (Bohm, 1980) with (e&eb) Carl Jung is theory of synchronicity (Jung and Pauli, 1955), with (e&eb) the implicate order providing (eb) a holographic medium through which apparently disconnected individuals become connected. The principle of synchronicity, the instantaneous connection of people and events beyond (e) the senses, has been equated with the quantum-physical principle of non-locality (Combs and Holland, 1990; Germine, 1991), and has been proposed to be the fundamental mechanism of conscious process (Germine, 1991).Applying the Holographic Principle theory of mind, subjects may be connected synchronously in the manner of two individuals conversing through cellular phones, where the radio signals are transmitted through (e&eb) a distant satellite. The difference here is that the radio signals travel at the speed of light, and thus the there is some miniscule time lapse between the sending and receiving ends. The Holographic Principle implies (eb) a much more distant and instantaneous communication pathway, with (e&eb) the analogue of the satellite being the holographic boundary of the Universe, as well as a much richer communication, potentially involving a transmission of thoughts and feelings. Our ego-consciousness seems to mask (e) the universal relatedness implied by (e) the Holographic Principle, and it is perhaps only through transcendence of (e) the ego-consciousness that the higher orders of experience can become conscious. In the phenomenon of synchronicity, there seems to be a meaningful connection between (e&eb) individuals that

breaks through the barrier of ego-consciousness. Such a connection is reported by many individuals in the course of dreams, when (e) the ego-consciousness has been suspended, at times informing (e&eb) the dreamer of something that has happened in the life of a meaningfully-connected individual. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Impassioned-intense Connectionist neuroscience, in an effort to attain relevance to clinical practice, has recently given us a theory of the relationship of the unconscious to (e&eb) consciousness which is (=) dualistic (Viamontes and Beitman, 2007), an ad hoc theory of conclusions based on (e) assumptions, the most erroneous of which are that representations of (e) information are information, and that such representations elaborate themselves through (e&eb) a process of representations of representations in neuronal circuits (Bechtel and Abrahamsen,1991). As theory, the connectionist models are lacking (e) heuristic value and **are (=) inimical to the progress** of (e) basic and clinical neuroscience. This is not to denigrate or minimize the importance of neuronal connections, but these connections are rather like the wires in (eb) a radio that plays a symphony. The radio, in and of itself, is (=) incapable of playing (e&eb) the symphony. The radio doesn't do not write the symphony, nor does it broadcast the Passionate-raging symphony over the airways. Yet, if we loosen one wire, the symphony is (=) not manifested through the radio. The brain is (=) not subject to (e) the same vulnerability, in that it has (e) parallel and redundant networks of processing, and, it is in this sense that the connectivity of the brain becomes (eb) important. Connectionism, to the extent that it accepts (e) consciousness and the (e&eb) unconscious at all, attributes them to (e) separate groups of "circuits" in the brain, with (e&eb) consciousness-processing circuits having (e) a limited but detailed ability to analyze (e&eb) information, and unconscious-processing processing more information in less detail (Viamontes and Beitman, 2007). There is an element of truth in this view in that information that has importance and meaning resonates within the brain at higher levels of experience and at higher levels of holographic recursion. This resonance involves what we know, what we feel, and what we are capable and willing to consciously realize. The Universe, as William Blake noted,

can be seen in a grain of sand. Indeed, the grain of sand requires (e) the Universe to exist. The human organism requires (e) the Self to exist. The Self is not a simple amalgam of (e) sensory experiences, as suggested in the connectionist model (Blinder, 2007). This amalgam view of the Self can be damaging when applied to (e&eb) psychotherapy, as the connectionists suggest (Blinder, 2007). The Self is ours forever and ever, from eternity to eternity. As Milton suggested, the mind is (=) its own place. Our deepest purpose is for our minds to resonate with (e&eb) the supra-conscious levels of experience, much as it is (=) purpose of the radio to play the symphony. If the symphony is not to our liking, (eb) we can turn the radio off. So, perhaps, the transpersonal consciousness can turn (e&eb) our own individual and collective radios off, or perhaps turn the volume down, if (e) we are not in harmony with the **Universal Mind**. Physics has given us many enigmas. Quantum theory itself is an enigma. But some of the enigmas seem to involve (e&eb) the limitation of (e) our own abilities of discernment. The thought experiment of (e) "Schrödinger's Cat," as it is often interpreted (e.g. Gribbin, 1984), assumes (eb) that the **cat is not conscious, and is (=) incapable of observation in the context of (e) collapse of (e) the wave function.** This is due, in large part, to the false dichotomy of (e) information and (e&eb) experience. Our observation of other animals indicates, to those that are attuned to (e) the animal mind, who can "feel" the animal mind, that, at (eb) the very least, all mammals are (=) conscious. It would seem that the great neuropsychiatrist, Stanley Cobb (1948), may have been correct in his attribution of (e) levels of consciousness to a wide variety of organisms: "...lower animals with (e&eb) no cerebrum appear to be conscious...even plants such as the sunflower that turns towards strong light may have (e) a vague awareness or warmth and comfort. There are many degrees of consciousness and it is my contention that it is integrated at (e&eb) many levels like other important functions of (e) the central nervous system."With respect to the theory of mind, recent mainstream thinking remains (eb) classical and mechanistic. Such views of mind are advocated by our most in prominent neuroscientists (e.g. Changeux, 1985; Gazzaniga, 1985) as neuroscience research attempts to explain all mental function on the basis of (e) **pure mechanism and the localized function of**

**specific areas of the brain.** Francis Crick, the distinguished Nobel Laureate, had a very influential second career as a neuroscientist. He concluded that we are basically soulless creatures, that all mental processes could be reduced to (e&eb) identifiable, neural correlates in (eb) the brain, and that, in fact, we are (=) nothing but a collection of neurons (Crick, 1995). Crick even went so far as to say that what we call the soul is (=) a group of neurons located in the prefrontal lobes of the brain. Another Nobel Laureate, Gerald Edelman (1987), has written a critically acclaimed book on "neural Darwinism," which purports that the brain is "circuitry" is constructed through a developmental process of "survival of the fittest" neurons and neuronal connections. At the same time the total purposelessness and lack of direction in evolution is expounded, again to critical acclaim, by Richard Dawkins (1986). The idea that there essentially no self or soul imbued with agency and self determination was expounded by Gilbert Ryle (1949/2002) with his parody of the "ghost in the machine," and this work continues to be influential among many scholars of the mind. Philosophers such as Patricia Churchland (1986), who coined the term "neurophilosophy," take an "eliminative materialist" view of (e) belief, free will, and consciousness. Other philosophers, such as Daniel Dennett (1991) would make (eb) experience "epiphenomenal," a by-product of (e) brain processes with (e&eb) no effect whatsoever on (eb) the function of mind. However, unlike neuroscience, which is still in its infancy, the philosophy of mind has had a long history, including such notable figures as Plato, Aristotle, Descartes, Kant, Locke, Hume, Kierkegaard, and James, to name just a few. Perhaps, as noted by Ervin Laszlo (1974): "in today's world, most of the traditional functions of (e) cognitive synthesis have atrophied (e) and are ignored and neglected." Hopefully, what we have presented here will be a step in the right direction for what Laszlo calls a "conceptual synthesis" to help fill the need for (e) meaningful engagement is such a world. The fundamental elements of such a synthesis have been described by (e) Laszlo as follows: Conceptual synthesis performs (e&eb) at least five basic functions in the guidance of affairs. They are the mystical, the cosmological, the sociological, the pedagogical or psychological, and the editorial functions. The mystical function inspires (eb)

in man a sense of mystery and profound meaning related to (e&eb) the universe and of himself in it. The cosmological function forms (eb) images of the universe in accord with (e&eb) local knowledge and experience, enabling (eb) men to describe and identify (e&eb) the structure of (e) the universe and the forces of nature. The sociological function validates supports and enforces (eb) social order, representing (eb) it in accord with (e&eb) the nature of the universe, or as the natural or right form of (e) social organization. The pedagogical or psychological function guides (eb) individuals through stages of life, teaching ways of understanding themselves and others and presenting (eb) desirable responses to life is challenges and trials. Finally, the editorial function of (e) conceptual is to define (eb)some aspects of reality as important and credible and hence to be attended to, and other aspects unworthy of serious attention. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Kant sums up this inversion and its spirit early in the Critique of Pure Reason: Up to now it has been assumed that all our cognition must conform to the objects; but all attempts to find something about them a priori through concepts that would extend our cognition have, on this presupposition, come to nothing. Hence let us once try whether we do not get farther with the problems of metaphysics by assuming that the objects must conform to our cognition, which would agree better with the requested possibility of an a priori cognition of them, which is to establish something about objects before they are given to us. This would be just like the first thoughts of Copernicus, who, when he did not make good progress in the explanation of the celestial motions if he assumed that the entire celestial host revolves around the observer, tried to see if he might not have greater success if he made the observer revolve and left the stars at rest. (B xvi) Onticology, like all variations of object-oriented ontology, is realist in its orientation. In defending a realist ontology onticology holds that the vast majorities of objects, actants, beings, or entities are independent of humans and are what they are regardless of whether any humans regard them or register them. In short, onticology rejects any anthropomorphic, idealist, or anti-realist thesis to the effect that to be is to be the correlate of mind, spirit, the body, the human, and language or otherwise. While it is certainly the case that knowledge is necessarily dependent on the object to which it relates, the converse does not hold true. Objects are not dependent on being known, regarded, perceived, or spoken

about. As such, and to put it in Aristotlean terms, knowledge is an accident of objects, not objects an accident of knowledge. As Althusser so nicely puts it, "[n]o doubt there is a relation between thought-about-the-real and thisreal, but it is a relation of knowledge, a relation of adequacy or inadequacy of knowledge, not a real relation, meaning by this a relation inscribed in that real of which the thought is the (adequate or inadequate) knowledge" (Reading Capital, 96). Althusser goes on to remark that "[t]he distinction between a relation of knowledge and a relation of the real is a fundamental one: if we did not respect it we should fall irreversibly into either speculative or empiricist idealism" (ibid.). Onticology categorically endorses Althusser's verdict. It is a fundamental necessity to distinguish between those relations that belong to the object and those that belong to knowledge. Contemporary philosophy continuously confuses these two very different sorts of relations. Naturally the question arises of how it is possible to surmount our relation to the object so as to determine whether objects themselves possess the properties we encounter in relating to objects. In other words, given that we can only ever relate to the object in relating to the object how is it possible to surmount this relation to get at the being of the object itself? Much more will have to be said about this later– and the answers will be surprising with respect to standard prejudices about realism –however, for the moment it can be said that onticology takes its epistemological inspiration from the transcendental realism of Roy Bhaskar. Among other things, Bhaskar sought to provide a transcendental grounding for the sciences. Insofar as onticology defends the thesis that the field of being is much more vast than the field of objects investigated by the natural sciences, it parts way with the thesis that the domain of being is exhausted by the domain of natural objects. However, the general form of Bhaskar's argument holds for our realist purposes. A transcendental argument seeks to elucidate the conditions under which certain acknowledged practices and forms of cognition are possible. Kant, for example, asked what must be the case for mathematical judgments to be possible. How is it both that we are able to extend our knowledge, as if by magic, through mathematical judgments and, more significantly, that these judgments are able to provide genuine knowledge of the world despite the fact that these

forms of reasoning are not based on experience? Part of Kant's argument consisted in claiming that mind imposes the forms of space and time on the data of experience. In other words, space and time are not attributes of being itself but rather of the mind that regards being. Insofar, Kant argues, as mathematics is ultimately a rumination on the nature of space and time taken in their most abstract form and insofar as the mind imposes space and time on the manifold of sensation, it thus follows that a priori judgments about the nature of spatio-temporal relations are possible that anticipate the structure of actual-space times without directly experiencing these space-times. Why? Because any manifold of sensation must necessarily be structured by these forms imposed by intuition. **Onticology–A Manifesto for Object-Oriented Ontology Part I Posted by larval subjects under Object-Oriented Philosophy My Question Is This: You Say, "Granted There Are Deeper And Deeper And More Primitive Forms Of "Knowledge" Or Abstraction Into Mechanism, But I Wouldn't Say That The Atomic Level Brings A New Epistemic World Into View." But Doesn't Any Kind Of Knowledge, However Primitive, Imply Bringing Some Kind Of Epistemic World Into View Or Proto-View?** Joel: David, Yes Exactly. My Question Is, Just How Far Down The Complexity Gradient Can The Meaning Of 'Knowledge,' Or 'Semiotics, Or Specifically Representation (Maya) Be Stretched Without Breaking? I Can Comfortably Stretch It To The Genetic Level, As The Proto-Epistemic, Where Genetic Codes 'Represent' Phenotypic Structures Or Modes, And The Whole Of Evolution Learns From Its Mistakes And Successes As That Knowledge And History Is Encoded Into And As The Recapitulation Of The Embryogenesis Of The Organism. But An Atom, In My Model, While Certainly Reacting To Its Environment From Its Own Attractor, Essence, And 'Interiority' (Prehension) Isn't Making Representations Of It To Do So. But The Key To The Power And 'Essence' Of Representation (And Simultaneously Its Limitation) Is The Capacity Of Abstraction And Choice. Evolution Does This At A Rudimentary Level (And Anthropomorphizing A Bit) With Its Sending Into The World Its 'Random' Variations, Organismic Ideas, In A Sense. Its 'Mind' Is Literally On The Outside As Our Living World And Biosphere. A 'Choice' Is Made In The Biospheric Interiority,

And Knowledge For The Evolutionary/Embryogenetic Trajectory Is Gained By The "Differential Reproductive Success" Of These Individuals Lives Through Time, As The Increasing Intelligence, Representation, And Capacity For Choice Is Directly Injected Into The Flow Through Sexual Selection. But Where Is The Encoded/Abstracted (Enfolded) Representation Of Options And A Choice Among Them Made At The Atomic Level? I Am Working With The Most Cutting-Edge And Coherent Models For The Atom (Including My Own) That I Can Find Among The Heterodoxy, And I See No Evidence Of Representation And Choice At This Level, Or Rather No Real Way To Further Stretch The Meaning Of The Terms To This Level. Sure There Is Uncertainty From Infinite Difference And Immanent Causation (See Bohm On Infinite Causation), But Without Representation There Really Is No Choice (Although Crudely We Could Call A 'Bifurcation' A 'Choice' By The System As A Whole). So I Am Simply Saying That Representation At The Atomic Level Is Entirely Enfolded And These Are Really Just Ontic Level Phenomena. The Atomic Level, In My View, Is Pre- Not Yet Proto-Epistemic, Which I Place at the Genetic Level Where We Find the First Codes And (e&eb) Primitive Representation. This Just Means That There Is A Real Distinction Between Ontic And Epistemic (And A Gradient Between Them) And We Can't Collapse The Two Into Just The Ontic-Epistemic. And Yet, With The Distinction Intact, And The Gradient To Explore, They Remain Nondual, Univocal, ONE Or Ontic. What I Would Say Exists At All Levels (My Own Self-Aware Myth About The Given) Is What I Call The "Symbiogenesis Of Subject And Object". This Is The Spinozan Essence Of Dynamic Stability And Growth, And Leibnizian Prehension. Also Called The "Nucleation Of Observability" In Spinbitz. It Is From This Nucleus Of The Observer (Fuller) In A "Point-Free Geometry" (Whitehead) Of Pointless "Points Of View" As Conditions Of Boundary That Wilber's "All Is Perspective" Finds Its Ontic Grounding. And This Symbiogenesis Is Also A Rudimentary Or Deep-Level Native Or Primitive Intelligence, Evolution At The Involution Of The Subject-Object Interface. It Is An Exploration Of Creative Learning And Direct Immanent Awareness, Unmediated (Enfolded) By Representational Forms (Unfolded). It Just Is, In Waves (Cosmic Vertebrae)

Before It Abstracts, Represents And Incarnates The Flesh Of What Could Be. Matter Is Enfolded Seed Of The Flower Of Its Abstraction. Maya Is This Depth And Recursion Of Brahma Into And Through The Boundary Conditions Of Abstract Relation. And Brahma Is The Span Of The Emergent Rigor Of The Cosmic Ergodic Spine That Opens The Recursions Into The Field Of Infinite Difference. In Immanent-Transcendent Waves, Literally And Empirically A Limit-Cycle In Ergodicity, It Enfolds And Unfolds From Anatom To Anatomy, From Infinite Intelligence Or Omnirelational Awareness Of Self, Through The Representational Depths Of Self As Self-Consciousness. Omni-Evolution, Trans-Dynamics, And The Anatom Posted On February 21, 2013A Discussion With Tom Huston, Part II The Following Is Part 2 Of An Excerpt From A Discussion Initiated By Tom Huston On His Facebook Page, Edited And Embellished For Clarity. Find Part One, Here. In This Part We Dig Into The Nature Of The "Anatom" And The "Ergodic Spine" Of An Electro-Fractal Cosmos, And Into The Trans-Dynamic Integration Between Being And Becoming, And The X-Interface−The Crossroads Of The Ontic-Epistemic (Brahma And Maya) And The Subject-Object Polarities. Univocal_Dynamics_V1 Tom Huston Is A Founding Member Of Integral Institute And A Former Editor Of Enlighten Next Magazine. His Webpage Is Tomhuston.Com. David Marshall Is A Writer And Editor Living In Chicago. No meaning, value, or normativity to be found in nature, that there is nothing **natural beings ought to be,** but that, rather, these judgments arise from us (See last post of larval subjects) which has led some to raise valuable questions about the coherence of these claims. The problem arises when the following three propositions are taken together: 1. there is nothing outside of nature.2. Beings have no intrinsic meaning, purpose, or value (in and of themselves, there's nothing they ought to be). 3. Value judgments about what beings and being ought to be arise from us and beings like us (bonobo apes, dolphins, institutions, birds of paradise, etc). The problem arises between proposition 1 and 3. How can it both be true that there is nothing outside of being and those normative judgments belong to us and other beings capable of making normative judgments, not nature? The problem arises from restricting these judgments to humans and beings capable of

making these judgments. In making such a claim it seems as if we're saying that there's something outside of nature, something that is beyond nature, thereby violating the first thesis and potentially reintroducing the nature/culture distinction. By contrast, if we say that normative judgments are the special domain of those living beings with the proper degree of sentience to make such judgments, then we seem to reintroduce the nature/culture distinction and fall back into the sorts of problems that I outlined in my last post (and that are so nicely critiqued by thinkers such as Latour). Is there a way out of this? I don't know. Great_Chain_of_Being_2_ (lighter) ASIDE: It seems to me that this is really what the debate between realism and anti-realism, realism and socio-linguistic constructivism surrounding the new materialisms and speculative realists has really been all about.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \qquad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9}+\frac{N_9 dG_9}{dR_9}-\frac{dB_9}{dR_9}>\frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9}+\frac{dM_9}{dR_9}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usance bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_9=\frac{N_9 G_9}{2M_9}\frac{d(9nF_9)}{dR_9}(\geq 0) \qquad (9a)$$

$$I_9=-\frac{N_9 G_9}{2M_9}\frac{d(9nM_9}{dR_9}(\geq 0) \qquad (9b)$$

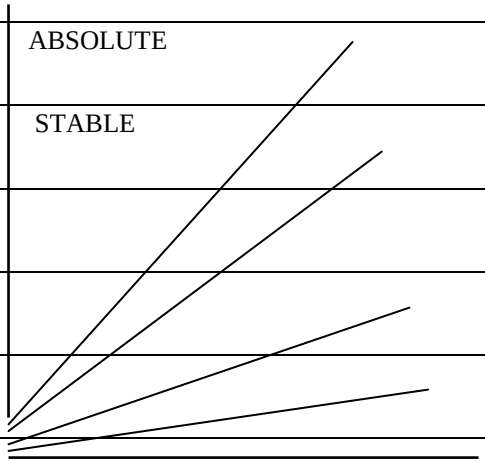
It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9>L_9-I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_9>L_9-I_9$, when $P_9^2/4>I$ (8a)

$P_9>L_9-I_9+\frac{(I_9-P_9^2}{(L_9-I)}$ when $P_9^2/4<I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Reffer couffesy Nunney

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability

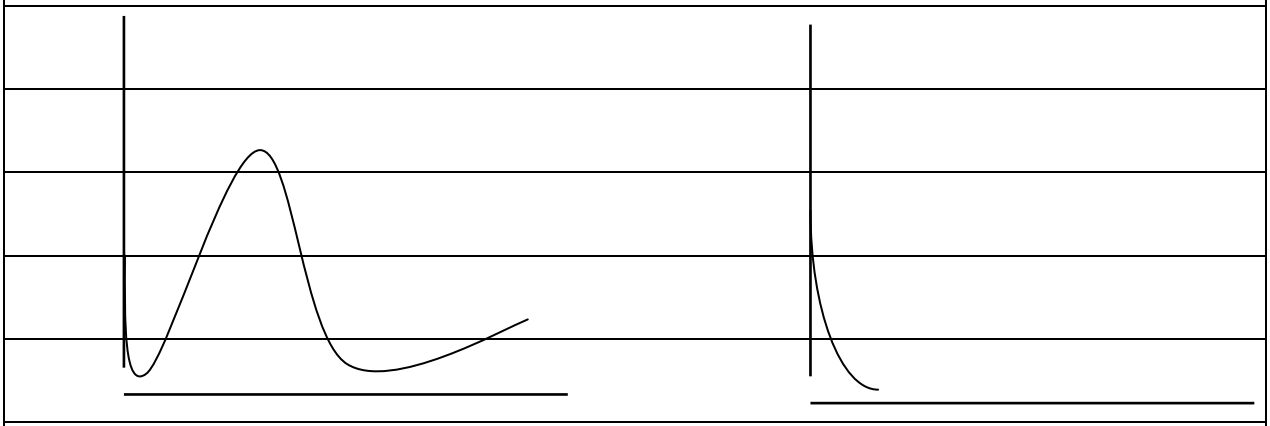
 

is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

Indeed, Nunny's model can be applied to some of the consciousness and philosophical problems. Predator prey analysis has ripped open the veneer of the fundamental concept of universe namely Debit-Credit world.

The noise of ordinary consciousness and (e&eb) beliefs drowns out (e) the signal. Unconsciousness is the background of (e) our ordinary awareness. Our organism is very much at (eb) the center of such effects. The organismic source is (=) our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is a meme-Scape. Pre-conceived concepts vie with structures, concepts with images. Like scientists who ignore assumed truths, we leapfrog over our beliefs and personality deficits, claiming idiosyncratic imagination is literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is both. We inherit the structures that make our experience what it is.

But the structure itself is “empty,” and each human culture “fills” it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us

feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through 're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns, nomenclature, and internal coherence of meaning and representation. In a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist

simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** The lower domain of experience is what we do not directly observe, the quantum realm, while the higher level is what we ordinarily observe, the classical realm. Based purely on experiment, the formulation of quantum theory initially placed the great divide between the quantum and classical domains between the measuring apparatus and the particle (Herbert, 1985). However, in the early 1930s, John von Neumann, in his rigorous mathematical treatise on quantum mechanics, found no support for such a division, compelling him to conclude that the wave for was collapsed by consciousness. Von Neumann's treatise has been called the "quantum bible" and "the most influential book on quantum theory ever written" (Herbert, 1985).

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t-T_9)F_9\big(R_9(t-T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t-T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) and the effective prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting $dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

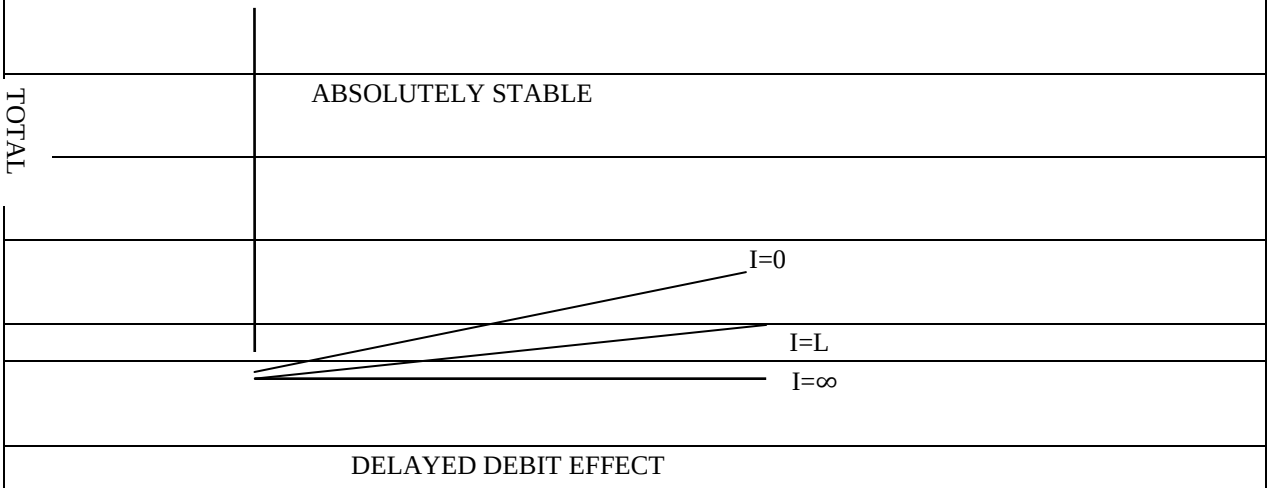



FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.

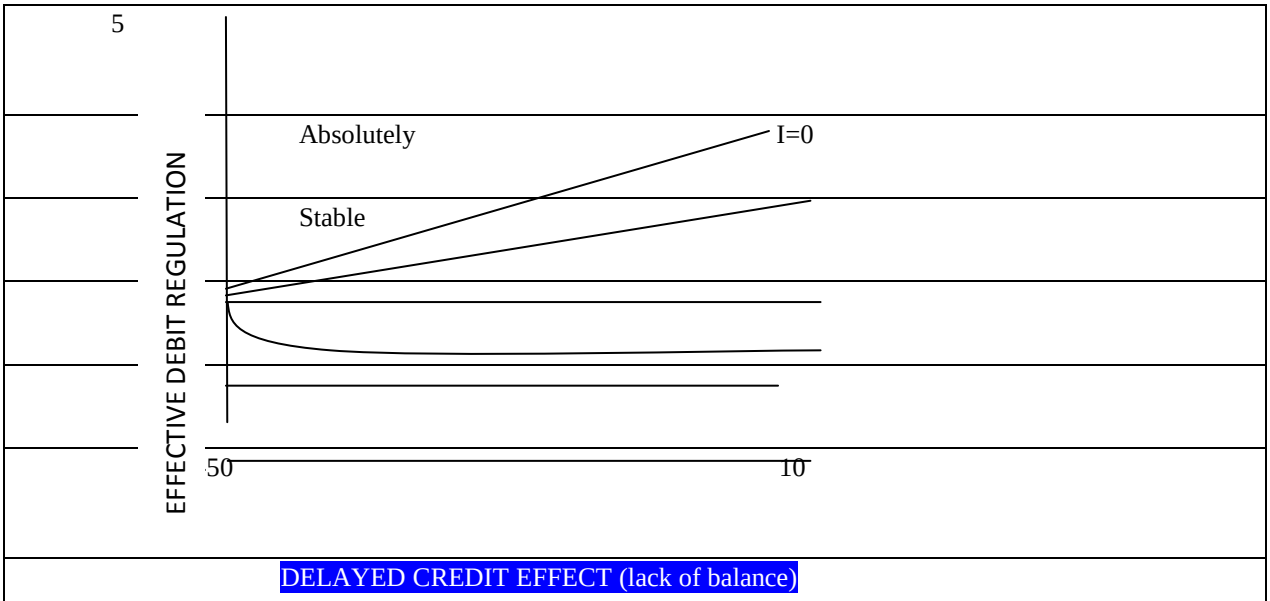



Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the

constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^\infty \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9\left(U_9 - V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9\right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$ then given the constraints of $f(T_9)$, $(A1.5)$ *can never be satisfied* Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDERATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 - V_9M_9 -]e^{-T_9\lambda_9} = 0$$

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9Z_9\sin(Z_9T_9) - (A_9M_9 + Y_9) + (A_9M_9 + X_9)\cos(Z_9T_9) = 0$$

$$(A_9M_9)Z_9 - M_9Z_9\cos(Z_9T_9) + (A_9M_9 - X_9)\sin(Z_9T_9) = 0$$

Where $A_9 = U_9 - V_9$ Using (A2.3), substitute for $Z_9^2$ and $Z_9$ in (A2.2):

$W_9^2(X_9^2 - 3M_9A_9X_9 - M_9^2(X_9 - Y_9) + 2M_9^3A_9 + 2A_9^2M_9^2]$

$$+W\left[MA^{3}-A^{2}\mathrm{X}+2\mathrm{MA}(2\mathrm{X}+\mathrm{Y})-2X^{2}\right]+A^{2}(\mathrm{X}+\mathrm{Y})=0$$

*where* $W_9$=1-cos ($\mathrm{Z}_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (\mathrm{X}_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ \mathrm{X}_9 - Y_9/2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range 0⍰$W_9$ ⍰2 *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (\mathrm{X}_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\ \mathrm{M}_9)^2\ (2A_9\mathrm{M}_9 - (\mathrm{X}_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9M_9^2 - 3M_9X_9\right) + W_9\left(3M_9A_9^2 - 2A_9X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9 + Y_9\right)$$

$=2M_9^2W_9^2(M_9 + A_9) + M_9A_9^2W_9 + \left(2A_9M_9 - X_9 + Y_9\right)(A_9W_9 + M_9W_9^2$

$+ (2- W_9)((2M_9X_9 + M_9Y_9)W_9 + A_9(X_9 + Y_9))$,

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range 0≤ $W_9$ ≤2.Defining

$Q_9(W_9, A_9) = a_9W_9^2 + c_9$ *where* $a_9$ $b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

**Harald Atmanspacher, Hans Primas** show that such an epistemic perspective can be consistently complemented by Einstein's ontically oriented position. **Time, Quantum and Information 2003, pp 301-321 Epistemic and Ontic Quantum Realities Harald Atmanspacher, Hans Primas Robert W. Spekkens** present a toy theory that is based on a simple principle: the number of questions about the physical state of a system that are answered must always be equal to the number that are unanswered in a state of maximal knowledge. Many quantum phenomena are found to have analogues within this toy theory. These include the noncommutativity of measurements, interference, the multiplicity of convex decompositions of a mixed state, the impossibility of discriminating nonorthogonal states, the impossibility of a universal state

inverter, the distinction between bipartite and tripartite entanglement, the monogamy of pure entanglement, no cloning, no broadcasting, remote steering, teleportation, entanglement swapping, dense coding, mutually unbiased bases, and many others. The diversity and quality of these analogies is taken as evidence for the view that quantum states are states of **incomplete knowledge** rather than states of reality. A consideration of the phenomena that the toy theory **fails to reproduce, notably, violations of Bell inequalities** and the existence of a Kochen-Specker theorem, provides clues for how to proceed with this research program.DOI: http://dx.doi.org/10.1103/PhysRevA.75.032110 **Evidence for the epistemic view of quantum states A toy theory Phys Rev A 75, 032110 – Published 19 March 2007 Robert W. Spekkens In such specialized and autonomous areas, it is to the eminent and erudite authors to interpret the contradistinctions and contradictions' in the dove tailing explaining. This I feel would be yeoman service in the creation of Nature's general ledger. M. S. Leifer and O. J. E. Maroney** examine the relationship between quantum contextuality (in both the standard Kochen-Specker sense and in the generalized sense proposed by Spekkens) and models of quantum theory in which the quantum state is maximally epistemic. We find that preparation noncontextual models must be maximally epistemic, and these in turn must be Kochen-Specker noncontextual. This implies that the Kochen-Specker theorem is sufficient to establish both the impossibility of maximally epistemic models and the impossibility of preparation noncontextual models. The implication from preparation noncontextual to maximally epistemic then also yields a proof of Bell's theorem from an Einstein-Podolsky-Rosen-like argument. DOI: http://dx.doi.org/10.1103/PhysRevLett.110.120401 Received 25 August 2012Published 20 March 2013 © 2013 American Physical Society **Maximally Epistemic Interpretations of the Quantum State and Contextuality Phys. Rev. Lett 110, 120401 – Published 20 March 2013 M. S. Leifer and O. J. E. Maroney.** Epistemological, onitic and ontological aspects particularly of quantum mechanics have invited

profound thinkers to contribute their mite to the growing literature on the subject, which includes the contributions from the subterranean realm and ceratoid dualism of philosophy like that Chaos theory of Deleuze, Derrida , Kierkegaard and others. **Andrei Khrennikov** shows that the Dirac-von Neumann formalism for quantum mechanics can be obtained as an approximation of classical statistical field theory. This approximation is based on the Taylor expansion (up to terms of the second order) of classical physical variables – maps f: Ω → R, where Ω is the infinite-dimensional Hilbert space. The space of classical statistical states consists of Gaussian measures ρ on Ω having zero mean value and dispersion σ2 (ρ) ≈ h. This viewpoint to the conventional quantum formalism gives the possibility to create generalized quantum formalisms based on expansions of classical physical variables in the Taylor series up to terms of nth order and considering statistical states ρ having dispersion σ2(ρ) = h n (for n = 2 we obtain the conventional quantum formalism). **December 2005, Volume 18, Issue 7, pp 637-650 Date: 28 Nov 2005 Generalizations of Quantum Mechanics Induced by Classical Statistical Field Theory Andrei Khrennikov** It is widely accepted that consciousness or, in other words, mental activity is in some way correlated to the behavior of the brain or, in other words, material brain activity. Since quantum theory is the most fundamental theory of matter that is currently available, it is a legitimate question to ask whether quantum theory can help us to understand consciousness. Several approaches answering this question affirmatively, proposed in recent decades, will be surveyed. It will be pointed out that they make different epistemological assumptions, refer to different neurophysiological levels of description, and adopt quantum theory in different ways. For each of the approaches discussed, these imply both problematic and promising features which will be indicated. **Discrete Dynamics in Nature and Society Volume 2004 (2004), Issue 1, Pages 51-73 http://dx.doi.org/10.1155/S102602260440106X Quantum theory and consciousness: an overview with selected examples Harald Atmanspacher Copyright © 2004 Hindawi Publishing**

**Corporation. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.** Another thought provoking and investigatory article is by Peter G. Lewis et al. Perhaps the quantum state represents information about reality, and not reality directly. Wave function collapse is then possibly no more mysterious than a Bayesian update of a probability distribution given new data. **Peter G. Lewis et al** consider models for quantum systems with measurement outcomes determined by an underlying physical state of the system but where several quantum states are consistent with a single underlying state—i.e., probability distributions for distinct quantum states overlap. Significantly, they demonstrate by example that additional assumptions are always necessary to rule out such a model. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.150404 **Distinct Quantum States Can Be Compatible with a Single State of Reality Phys. Rev. Lett. 109, 150404 – Published 9 October 2012 Peter G. Lewis et al** Communication complexity of a quantum channel is the minimal amount of classical communication required for classically simulating a process of state preparation, transmission through the channel and subsequent measurement. It establishes and perpetuates a limit on the power of quantum communication in terms of classical resources. **Alberto Montina** shows that classical simulations employing a finite amount of communication can be derived from a special class of hidden variable theories where quantum states represent statistical knowledge about the classical state and not an element of reality. This special class has attracted strong interest very recently. The communication cost of each derived simulation is given by the mutual information between the quantum state and the classical state of the parent hidden variable theory. Finally, find that the communication complexity for single qubits is smaller than 1.28 bits. The previous known upper bound was 1.85 bits. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.110501 **Epistemic View of Quantum States and**

**Communication Complexity of Quantum Channels Phys. Rev. Lett 109, 110501 – Published 12 September 2012 Alberto Montina for models see elsewhere .** Stability analysis of what? Quantum state representing reality or our knowledge of reality? Does the quantum state represent reality or our knowledge of reality? In making this distinction precise, **Nicholas Harrigan, Robert W. Spekkens** are led to a novel classification of hidden variable models of quantum theory. They show that representatives of each class can be found among existing constructions for two-dimensional Hilbert spaces. Approach also provides a fruitful new perspective on arguments for the nonlocality and incompleteness of quantum theory. Specifically, **Nicholas Harrigan, Robert W. Spekkens** show that for models wherein the quantum state has the status of something real, the failure of locality can be established through an argument considerably more straightforward than Bell's theorem. The historical significance of this result becomes evident when one recognizes that the same reasoning is present in Einstein's preferred argument for incompleteness, which dates back to 1935. This fact suggests that Einstein was seeking not just any completion of quantum theory, but one wherein quantum states are solely representative of our knowledge. Hypothesis is supported by an analysis of Einstein's attempts to clarify his views on quantum theory and the circumstance of his otherwise puzzling abandonment of an even simpler argument for incompleteness from 1927. **February 2010, Volume 40, Issue 2, pp 125-157 Date: 09 Jan 2010 Einstein, Incompleteness, and the Epistemic View of Quantum States Nicholas Harrigan, Robert W. Spekkens** statistical model of the probabilistic description of physical reality has been studied by many authors. **Andrei Khrennikov posits** a contextual statistical model of the probabilistic description of physical reality. Here contexts (complexes of physical conditions) are considered as basic elements of reality. There is discussed the relation with QM. He proposes a realistic analogue of Bohr's principle of complementarity. In the opposite of the Bohr's principle, **Andrei Khrennikov** principle has no direct relation with mutual exclusivity for

observables. To distinguish his principle from the Bohr's principle and to give better characterization, he changes the terminology and speaks about supplementarity, instead of complementarity. Supplementarity is based on the interference of probabilities. It has quantitative expression trough a coefficient which can be easily calculated from experimental statistical data. There is need not appeal to the Hilbert space formalism and noncommutativity of operators representing observables. Moreover, in our model there exists a pair of supplementary observables which cannot be represented in the complex Hilbert space. There are discussed applications of the principle of supplementarity outside quantum physics. **October 2005, Volume 35, Issue 10, pp 1655-1693 The Principle of Supplementarity: A Contextual Probabilistic Viewpoint to Complementarity, the Interference of Probabilities and Incompatibility of Variables in Quantum Mechanics Andrei Khrennikov Foundations of Physics** Description of the vacuum in Yang-Mills theory and arrive at a physical interpretation of the pseudoparticle solution and the attendant violation of symmetries has been done by **R. Jackiw and C. Rebbi.** It is here vaccuum stability forms the bastion, pillar, post and stylobate of the mechanisms of yang mills Theory and action of a particle. The existence of topologically inequivalent classical gauge fields gives rise to a family of quantum mechanical vacua, parametrized by a CP-nonconserving angle. The requirement of **vacuum stability** against gauge transformations renders the vacua chirally noninvariant. DOI: http://dx.doi.org/10.1103/PhysRevLett.37.172 **Vacuum Periodicity in a Yang-Mills Quantum Theory Phys. Rev. Lett 37, 172 – Published 19 July 1976 R. Jackiw and C. Rebbi.** Three-dimensional Yang-Mills and gravity theories augmented by gauge-invariant mass terms are analyzed. These topologically nontrivial additions profoundly alter the particle content of the models and lead to (eb) quantization of a dimensionless mass-coupling-constant ratio. The vector field excitations become (=) massive, with (e&eb) spin 1 (rather than massless with spin 0), and the mass provides (eb) an infrared cutoff. The gravitation

acquires (e) mass, mediates (e&eb) finite-range interactions, and has (e) spin 2 (rather than being absent altogether); although its mass term is of (e)third derivative order, there are no ghosts or acausalities. DOI: http://dx.doi.org/10.1103/PhysRevLett.48.975 © 1982 The American Physical Society **Three-Dimensional Massive Gauge Theories Phys. Rev. Lett 48, 975 – Published 12 April 1982 S. Deser et al** All scattering amplitudes in (eb) the maximally supersymmetric N=4 super-Yang–Mills theory possess (e) a new, dual superconformal symmetry which extends the previously discovered dual conformal symmetry of MHV amplitudes. To reveal this property authors formulate the scattering amplitudes as (=) functions on (e&eb) the appropriate dual superspace. Rewritten in this form, all tree-level MHV and next-to-MHV amplitudes exhibit manifest dual superconformal symmetry. Authors propose a new, compact and Lorentz covariant formula for (e) the tree-level NMHV amplitudes for (e) arbitrary numbers and types of external particles. The dual superconformal symmetry is broken at loop level by infrared divergences. However, authors provide evidence that the dual conformal anomaly of (e) the MHV and NMHV superamplitudes is (=) the same and, therefore, their ratio is (=) dual conformally invariant. Authors show this explicitly for (e) the six-particle amplitudes at one loop. Authors conjecture that these properties hold for all, MHV and non-MHV, superamplitudes in N=4 SYM both at (eb) weak and at strong coupling. **Nuclear Physics B Volume 828, Issues 1–2, 21 March 2010, Pages 317–374 Dual superconformal symmetry of scattering amplitudes in N=4 super-Yang–Mills theory J.M. Drummonda DOI: 10.1016/j.nuclphysb.2009.11.022.** By comparison with numerical results in (eb) the maximal Abelian projection of lattice Yang–Mills theory, it is argued that the nonperturbative dynamics of (e) Yang Mills theory can be described by (e) a set of fields that take their values in (eb) the coset space SU(2)/U(1). The Yang–Mills connection is parameterized in (eb) a special way to separate (e&eb) the dependence on the coset field. The coset field is then regarded as (=) a collective variable and a method to obtain its effective

action is developed. It is argued that the physical excitations of (e) the effective action may be (=) knot solitons. A procedure to calculate the mass scale of knot solitons is discussed for (e) lattice gauge theories in the maximal Abelian projection. The approach is extended to (e&eb) the SU (N) Yang–Mills theory. A relation between the large N limit and the monopole dominance is pointed out. **Physics Letters B Volume 458, Issues 2–3, 8 July 1999, Pages 322–330 An effective action for monopoles and knot solitons in Yang–Mills theory Sergei V. Shabanov DOI: 10.1016/S0370-2693(99)00612-7.** It is here the quintessential formulation of consciousness occurs. Like a man standing on the threshold of infinity trying to ponder what lies beyond the veil, which separates the seen from unseen a true seeker tries to lift the individual general ledger to the cosmic general ledger. Like in Darwinian evolution, mind also has to evolve from individual consciousness to cosmic consciousness. In Kaishmiri Shaivism one withdraws from the finite to the infinite, but one also goes on an outward journey from the infinite to the finite, because both the finite and the infinite have an intimate connection. The finite is not seen as unreal, but as a symbol of the infinite. There is no real distinction between them. Those two movements constitute Spanda, a key concept in Kaishmiri Shaivism. Spanda is the pulsation of the Absolute in different phases of being. There are no opposites like subject and object, unity or duality, absolute or relative. They are just different phases of the universal vibration of the Absolute. The goal is to realize or be (=) at once infinite and finite. One does not turn away from appearances (like in Advaita Vedanta), but one realizes that the Absolute manifests all things. Spanda, the eternal pulsation of the Absolute, oscillates between a passion to create and dispassion from the created. Through it the Absolute transform itself into all things and then returns back into the emptiness of its undifferentiated nature. In Kaishmiri Shaivism the Absolute is seen as pure consciousness (=being). The Absolute is an eternal all-pervasive principle, the highest reality, the nature of all entities eternally and blissfully at rest within its own nature. The Absolute is the nature of the Self (and

thus of us all). The Absolute is divine; it is Shiva, the Lord of the Universe. It is full of (e) conscious activity through (e&eb)which it generates the universe, and reabsorbs it into (e&eb) itself at the end of each cycle of creation. Thus we speak of monism, as everything resides within (eb) this one absolute consciousness. It sustains (eb) all things, it embraces (e) all things, and it pervades (e&eb) all things. All things are appearances within (eb) the absolute consciousness, but nevertheless real (in contrast to (e) Advaita Vedanta where (e) appearances are seen as unreal or illusionary). All things appear external (out there, outside ourselves), but they do not have (e) a being on their own. They do not exist as (=) separate entities on their own. Everything is contained within (eb) consciousness. What we see as objects are (=) manifestations of consciousness. The event which constitutes (eb) the universe is always internal events happening within (e&eb) consciousness because (e) their essential nature is consciousness itself. If a physical object were totally material, and independent or external to consciousness, it could never be experienced. The universe and (e&eb) consciousness are two aspects of a whole. The universe is (=) an attribute of consciousness which bears consciousness as (=) its substance. Doctrine of Vibration: S.G.Dyczkowski what you "see" is (=) not always what you "get." Many people mistakenly take their own visions literally without (e) "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in (eb) the physical sense but in a psychological sense. The archetypes of (e) the collective unconscious are arrayed behind (e) the scenes of current worldwide conditions, of (e) crisis and confusion. They mirror (e&eb) our own states back at us, whether we perceive them as (=) such or interpret them plausibly or not (Note the essay has been written in model able form and models are given in various papers.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4\,a_9 \square - b_9 \square 0$$

And

$$b_9^{2} \square 4\,a_9 c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9\, in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\, X_9$ ,

$Let\ Y_9 = \ y_9\, X_9$And try the substitution $A_9 = X_9/(2 + A_9\ y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + \ n_9 y_9)$ we have

$$W_9^{2}\, y_9 \left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^{2}\left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9\, X_9 \left[\frac{X_9 (1 + n_9 y_9)}{M_9^{2}(2 + n_9 y_9)^{2}} + 2\, y_9 \left(n_9 - 1\right) + \frac{X_9^{2}(9 + y_9)}{M_9^{2}(2 + n_9 y_9)} = 0.\ (A2.12)\right.$$

Since $n_9$ and $A_9$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9 = (2 + n_9 y_9)$, *the condition* (A2.99?) *can be written as*

$[\frac{X_9(E_9-9)}{E_9^2}]^2[\frac{X_9^2}{M_9^2} - 4\, y_9\ E_9^2 \frac{X_9}{M_9^2} + 4y_9 (n_9 - 2)E_9^3]$ .

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9 / W_9^2)$ can be expressed as

$2\, E_9 [y_9 E_9 + (y_9\ E_9 (9 + y_9)(E_9 - n_9))^{9/2}] > \frac{X_9}{M_9^2} > 2E_9 [y_9 E_9 - y_9 E_9 (9 + y_9)(E_9 - n_9))^{9/2}]$

Concept Of Being From Shaivite Point Of View: In His Famous Work Entitled On The Conservation Of Force (1847), Hermann Von Helmholtz Had Divided The Totality Of The Energy In The Universe Between Potential Energy And Kinetic Energy And Affirmed The Reciprocal Convertibility Of The Two. In 1852, William Thomson Pointed Out That There Exists A Sub-Ensemble Within Kinetic Energy, Heat, Which, Once It Has Been Generated, Is No Longer Entirely Convertible Into Potential Energy-Or Into Any Other Form Of Kinetic Energy. Considering That The (Partial) Reconversion Of Heat Into Labor Is Possible Only In Situations That Present A Disparity In Temperature, And That Heat Tends To Pass From Warmer To Cooler Bodies By Spreading On An Even Temperature Level Through Space, Thomson Concluded That The Universe Tends Towards A Final State Where Any Energetic Transformations, Every Movement And Every Form Of Life Will Cease: The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO: The Ultimate Nature Of Anything Cannot Change. Change Of Any Sort Is Merely Apparent (Vivarta). The World Of Change And Becoming Is A False Super- Imposition (Adhyaropa, Adhyasa) On The Absolute. It Appears So On The Screen Of Individual, Collective Consciousness, Albeit One Cannot Be So Sure Of Cosmic, For There Are Many Reasons That Cosmic Transactions Also Had To Be Subjected To Change

By ParaBrahaman. In Cosmic Terms, The Mistake (Bhranti) Consists Of The Supposition That The Real Brahman Is The Unreal Universe And The Unreal Universe Is The Real Brahman. In Microcosmic Terms, It Is The Mistake Of Falsely Conceiving The Body, Mind Or Even One's Personality To Be The Self. Note We Have Taken Self As Witness Consciousness. Notwithstanding There Is Nothing Wrong In The Conception Of Self As A Dynamical Presupposition, In That Individual Consciousness And Collective Consciousness And Even Cosmic Consciousness Evolutes In To ParaBrahaman. In The Same Way As The Image Of A Snake Is Falsely Superimposed On A Rope, Similarly The Universe Is Falsely Projected Onto The Real Substratum, The Brahman. Ignorance Is Not Merely A Personal Lack Of Knowledge, But A Cosmic Principle. As Such It Is Called "Maya," The Indefinable Factor (Anirvacaniya) That Brings This Mistake In Identity About. The Reality Status Of This Cosmic Illusion Is Also Indefinable: On The One Hand It Is Not Brahman, The Sole Reality; On The Other Hand It Is Not Absolutely Non-Existent Like A Hare's Horn Or The Son Of A Barren Woman. Note Here The Vedantin Avoids Talking About Zero Which We Have Avidly And Assiduously Stressed About. Nothing Is The Factor That Exists. Also See the Remarks about Satre's Being And Nothingness. Brahman Is The Source Of World Appearances Only In The Sense Of Being Their Unconditioned Ground Or Essential Nature. The Universe Is False Not Because It Has No Nature Of Its Own But Because It Does Have One. The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski. Concept Of Being From Shaivite Point Of View: It Is The 'I Am The Body/Mind' Belief That Gives Rise To The 'I Am Not The World' Belief. These Two Beliefs Are (=) Co-Created."Consciousness Projects The Appearance Of The Mind, Body And World By Taking (Eb) The Shape Of Thinking, Sensing And Perceiving."“Attention Is Consciousness With An Object. When The Object Vanishes, Attention Simply Remains What It Always Is Consciousness."“There Is No Purpose To Meditation. The Purpose Is Already Accomplished."“Everything

That Is Experienced Is Experienced By, Through, In And As Consciousness.”“The Seen Cannot Be Separated From Seeing And Seeing Cannot Be Separated From Consciousness.” “**The Reality Of Any Experience Is Not Hidden In The Appearance; It Is Expressed By The Appearance The Transparency Of Things**' By Rupert Spira Book Review By Dennis Waite The Following Is A Review Of Rupert's Book: 'The Transparency Of Things: Contemplating The Nature Of Experience’. There Is An Essay From The Book Here And The Essay On Cézanne May Be Read At The Awakened Eye Website Nature Rejects The Naiveté That Seeks Absolute Truth. We Are Beginning To Realize, Individually And Culturally, That "Realities" Are All Human Constructions. The Task Becomes One Of "Catching Ourselves In The Act" Of Creating Our Own "Reality" From The Flow Of Events. Human Truth Is Always An Engagement Of Mind With Experience. The Challenge Of The Therapist In These Times Of Chaotic Change Is To Validate The Concept That We Don't Have To Fear The Collapse Of What We Think We Are. Strange Attractors: Transference, Holography, And An **Archetype Burke, J. (2003). Strange Attractors: Transference, Holography, And An Archetype (Doctoral Dissertation, Pacifica Graduate Institute, 2003).** When we get to the topic and understanding of the Transcendental Logic, the brilliant and baffling Transcendental Deduction (TD) Recall the two movements just discussed, the one from experience to its conditions and the one from the forms of valid inference to the concepts that we must use in all judging (the Categories). This duality led Kant to his famous question of right (quid juris) (A84=B116): with what right do we apply the Categories, which are not acquired from experience, to the contents of experience? (A85=B117). Kant's problem here is not as arcane as it might seem. It reflects an important question: How is it that the world as we experience it conforms to our logic? In briefest form, Kant thought that the trick to showing how it is possible for the Categories to apply to experience is to show that it is necessary that they apply (A97).TD has two sides, though Kant never treats them separately John Marlado,

accounts, attributes, and ascribes the philosophy of Nishida KitarÅ is lucid, insightful, and deeply informative despite the highly questionable argument that gives the book its structure. In comparison with other, technically more accurate accounts of Nishida's philosophy, it is broader in scope than most but more single-minded in its approach. The analysis is for the most part limited to four books that have been translated into English, out of approximately 16 volumes of philosophical essays in Japanese, but it aims to capture the contours and most important details of Nishida's entire philosophy. To appreciate the achievement of this account, one would need only to read, in English translation or in the Japanese original, any of the volumes with their streams of repetitive, puzzling expressions and circuitous, often dead- end argumentation. Wilkinson's elucidation of details in the passages he samples makes it easier to understand why Nishida is considered Japan's greatest academic philosopher. A critical look at the frame of the book and some of its own claims reveals where, to my mind, it has gone astray. The argument that frames this book is repeated throughout: Nishida's philosophy is an attempt to articulate, in a characteristically Western philosophical manner, the Zen experience he had as a young man. Nishida first made the attempt by using the conceptual frameworks of Western philosophers, and after they failed him he developed his own, largely by recasting Buddhist insights. In the end, a philosophical articulation of the Zen worldview in Western terms or even in Eastern terms recast by conceptual, Western philosophical methods approved impossible. What Nishida attempted is incommensurable with Western philosophy. While the author first offers these statements as hypotheses (p. 2), one premise of his argument that there is such a thing as "the Zen world-view" or "the Zen experience “remains unquestioned; and another that Nishida's pivotal experience was the Zen experienceâis asserted as "beyond question" (p. 151; see also pp. 28, 48, 10, 114).Wilkinson's presentation of Zen in the first chapter is drawn largely from D. T. Suzuki, with help from Suzuki's portage Abe Masao; some comments on time, birth, and death in; and a few other sources. The

result is a caricature that imagines Zen as an invariable, unitary experience. In fact, Dagan's teachings on practice as the manifestation of enlightenment do not sit well with D. T. Suzuki's emphasis on a satori experience, and Suzuki does not enlighten us about what Zen in its various practices has historically been. The author would have benefited from acquaintance with other presentations by Zen teachers, such as that of ShunryÅ Suzuki'sZen Mind Beginner's Mind, with historical accounts of Zen traditions, and with contrasts between DÅgen's view of time and Sam Â· sÄraand that of other Buddhist philosophers. A little research in contemporary Zen scholarship would have shown how implausible is the talk of "nirvana or awareness of mu" (p. 24), "the basic Zen position [that] never varies" (p. 3), "the Zen vision" (pp. 5, 7), "Zen epistemology" , "Zen experience" (p. 59), "the conception of time involved in Zen" (p. 15), and "the fundamental Zen assertion that there is no absolute time" (p. 16). Even if DÅgen's notion that time and being are nondual (p. 16) epitomized most of Zen, Nishida's early notion of "a transcendent, unchanging reality apart from time" (p. 51) would contrast sharply with it. Equally questionable is Wilkinson's contention that Nishida's "pure experience" is equivalent to Zen satori. The notion of pure experience Nishida developed in his first major work may be pivotal for his entire philosophy, but he never claims or implies that any part of his philosophy presupposes an exceptional experience like satori that grounds his convictions and renders them unverifiable to those who have not had the experience, as Wilkinson intimates (pp. 55, 159, and 161). Nishida presents his notions of pure experience and intellectual intuition, the grasp of its unity, and later... Perennialism and Constructivism is studied by Randolph T Dible : experiencer or the observer and the transcendental subjectivity itself Pure objectivity is essentially a deconstruction, apophatic and negativa The encountering of the experiencer or observer—transcendental subjectivity itself—at the foundation of the world leads inevitably to the recognition of pure objectivity as ultimate reality (which can be taken as its ultimate deconstruction, analogous to the apophatic or via

negativa), from which objects derive their value, weight, significance, meaning or objectivity. In this way, pure objectivity can be seen as the supra-self-evident Axiological Axiom, so to speak, even Unconditional Love, in romantic terms. This axiology (value theory) has a structure inverse to the relationship between transcendental subjectivity as the radical unity of pure self-reference and on the other hand, the world of forms, as mere traces (representations, indications) of the unique, original "first distinction" Spencer-Brown speaks of at the foundation of his calculus. That is, all forms (i.e., distinctions, differences) would reduce to being the first distinction, also known as the marked state, which can be called penultimate reality (pure self-reference or transcendental subjectivity: the Spirit which animates us), except that forms are complimentary to their content, which is their objectivity or value, which would reduce to the unmarked state or ultimate reality. It is the incongruity of form (thoughts; Whitehead's "negative prehensions") and value (feelings; Whitehead's "positive prehensions," or my notion of objectivity, meaning and qualia; in short, the non-formal aspects of experience) that holds forms open and keeps them from absolute reduction. This accounts for the brute, concrete persistence of the "functional illusion"-- to use a term from Dzogchen Buddhism-- of the world. Thus this system has an axiology of metaphysical objectivity grounded on the ideal of pure objectivity as the source of all value, meaning and significance, itself the very fecundity of profundity, which is the motive of drawing the distinction in the first place **(Wikipedia, Kant's writings and Logic Of Sense)** Nano rings were created by accident while intending to make quantum dots. They have interesting optical properties associated with excitons and the Aharonov–Bohm effect. Application of these rings used as light capacitors or a buffer includes photonic computing and communications technology. Analysis and measurement of geometric phases in mesoscopic rings is ongoing. Several experiments, including some reported in 2012, show Aharonov-Bohm oscillations in charge density wave (CDW) current versus magnetic flux, of dominant period h/2e through CDW rings up to 85 **μm** in

circumference above 77 K. This behavior is similar to that of the superconducting quantum interference devices (seeSQUID). **Aharonov–Bohm nano rings (Wikipedia) Stability analysis of nano rings would help progressively accentuate the research about the materials division that uses nano rings.** From the profound contemplation and wisdom of the Buddha and Mahabodhisattvas, we know there is no such reality. Instead, egolessness (non-self) is the only path to understand the reality of the deluded life. All existences are subject to the law of causes and conditions. These include the smallest particles, the relationship between the particles, the planets, and the relationship between them, up to and including the whole universe! From the smallest particles to the biggest matter, there exists no absolute independent identity. Egolessness (non-self) implies the void characteristics of all existence. Egolessness (non-self) signifies (eb) the non-existence of permanent identity for self and existence (dharma). Sunyata stresses (eb) the voidness characteristic of (e) self and existence (dharma). Sunyata and Egolessness possess (e) similar attributes. As we have discussed before, we can observe the profound significance of Sunyata from (e) the perspective of inter-dependent relationships. Considering dharma-nature and (e&eb) the condition of nirvana, all existences are immaterial and of (e) a void-nature. Then we see each existence as (=) independent of each other. But then we cannot find any material that does exist independent of (e) everything else. So Egolessness also implies (eb) void-nature! From the observation of all existences, we can infer (eb) the theory of nirvana and the complete cessation of (e) all phenomena. From the viewpoint of phenomena, all existences are so different from (e) each other, that they may contradict (e&eb) each other. They are (=) so chaotic. In reality, their existence is (=) illusionary and arises from (e) conditional causation. They seem to exist on one hand, and yet do not (e) exist on the other. They seem to be united, but yet they are so different to (e) one another. They seem to exist and yet they do cease! Ultimately everything will return to (e) harmony and complete calmness. This is the nature of all existence. It is the

final resting place for all. If we can understand this reality and remove our illusions, we can find this state of harmony and complete calmness. **Teachings in Chinese Buddhism (6) sunyata (emptiness) in the Mahayana context (Wikipedia) Many models have been given concatenating Buddhist doctrines and Vaccum energy. These models essentially throw light on the essence of self(witness consciousness) and the cosmic general ledger (cosmic consciousness) while the imbalance in the gratification deprivation complex always remain unbalanced when it comes to individual general ledger , advocating the evolutionary theory of individual consciousness to cosmic consciousness. It is to be noted that it is not the question of knowing all the transactions but the wisdom of what lies behind the debit- Credit, Debit-Debit, or Credit-Credit principle that is most important for assimilation.** Look at this simple thesis to prove the point stated in the foregoing. Thesis: 'The World Has A Beginning In Time, And In Space It Is Also Enclosed In Boundaries.' Proof: 'For If One Assumes That The World Has No Beginning In Time, Then Up To Every Given Moment In Time An Eternity Is Elapsed, And Hence An Infinite Series Of States Of Things In The World, Each Following Another, Has Passed Away. But Now The Infinity Of A Series Consists Precisely In The Fact That It Can Never Be Completed Through A Successive Synthesis. Therefore An Infinitely Elapsed World-Series Is Impossible, So A Beginning Of The World Is A Necessary Condition Of Its Existence, Which Was The First Point To Be Proved (Kant, Critique Of Pure Reason, Tr. By R. Guyer And A. W. Wood (Cambridge: Cambridge University Press, 1998): B 454, P. 470). As Quoted In The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO"The Cyclical Hypothesis, So Heavily Criticized By Nietzsche (VP II 325 And 334), Arises In This Way." In Fact, Nietzsche Was Not Criticizing The Cyclical Hypothesis But Only The Particular Form Of That Hypothesis Presented In Vogt's Work. All Of Nietzsche's Texts Without Exception Speak Of The Eternal Return As The Repetitio Of The Same Events Within A Cycle Which Repeats Itself

Eternally. If Deleuze's Interpretation Holds That The Eternal Return Is Not A Circle, Then What Is It? A Wheel Moving Centrifugally, Operating A "Creative Selection," "Nietzsche's Secret Is That The Eternal Return Is Selective" Says Deleuze: The Eternal Return Produces Becoming-Active. It Is Sufficient To Relate The Will To Nothingness To The Eternal Return In Order To Realize That Reactive Forces Do Not Return. However Far They Go, However Deep The Becoming-Reactive Of Forces, Reactive Forces Will Not Return. The Small, Petty, Reactive Man Will Not Return. Affirmation Alone Returns, This That Can Be Affirmed Alone Returns, Joy Alone Returns. Everything That Can Be Denied, Everything That Is Negation, Is Expelled Due To The Very Movement Of The Eternal Return. We Were Entitled To Dread That The Combinations Of Nihilism And Reactivity Would Eternally Return Too. The Eternal Return Must Be Compared To A Wheel; Yet, The Movement Of The Wheel Is Endowed With Centrifugal Powers That Drive Away The Entire Negative. Because Being Imposes Itself On Becoming, It Expels From Itself Everything That Contradicts Affirmation, All Forms Of Nihilism And Reactivity: Bad Conscience, Ressentiment..., We Shall Witness Them Only Once. [...] The Eternal Return Is The Repetition, But The Repetition That Selects, The Repetition That Saves. Here Is The Marvelous Secret Of A Selective And Liberating Repetition. There Is No Need To Remind The Reader That Neither The Image Of A Centrifugal Movement Nor The Concept Of A Negativity-Rejecting Repetition Appears Anywhere In Nietzsche's Writings, And Indeed Deleuze Does Not Refer To Any Text In Support Of This Interpretation. Further, One Could Highlight That Nietzsche Never Formulates The Opposition Between Active And Reactive Forces, Which Constitutes The Broader Framework Of Deleuze's Interpretation. For Some Years, Marco Brusotti Has Called Attention To The Fact That Deleuze Introduced A Dualism That Does Not Exist In Nietzsche's Writings. To Be Sure, The German Philosopher Describes A Certain Number Of "Reactive" Phenomena (For Example, In The Second Essay Of The Genealogy Of Morality, § 11, He Talks About

"Reactive Affects" [Reaktive Affekte], "Reactive Feelings" [Reaktive Gefühlen], Reactive Men [Reaktive Menschen]); But These Are Nonetheless The Result Of Complex Ensembles Of Configurations Of Centers Of Forces That Remain In Themselves Active. Neither The Word Nor The Concept Of "Reactive Forces" Ever Appears In Nietzsche's Philosophy. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO Some Parts Are Deleted Due To Spatial Constraints. In the midst of thoughts about the eternal return we find at least two other thematic axes On the one hand, the view of the **world as a constant flux of forces without any goal, law, or rules of becoming. A chaos sive natura de-divinized and de-anthropomorphized which constitutes the "ontological substratum**" of the whole of Nietzsche's reflexions On the other hand, an ensemble of fragments of an anthropologico-sociological character, designing a path of liberation leading to the creation of superior individuals by way of a profound transformation of their instinctual structure. This transformation must be achieved by a practice of solitude and internal struggle towards the liberation from the ancient representations of the world and from the incorporated herd values. For an analysis of these thematic perspectives, see Paolo D'Iorio, La linea e il circolo. Cosmologia e filosofia dell'eterno ritorno in Nietzsche (Genova: Pantograf, 1995): 233-322. With most of the parameters quantified, with assumptions there must be picturesque view of the self and non self, space and time, science and philosophy with the concatenational consummations. After The Chapter "On Redemption," Where Zarathustra Dares Not Expose His Doctrine, The Eternal Return Begins To Be Enunciated In Part Three Of The Work. In The First Place, It Is The Dwarf Who Formulates It In The Chapter "On The Vision And The Riddle." Facing The "Gate Of The Instant" Which Symbolizes The Two Infinities That Stretch Towards The Past And The Future, The Dwarf Whispers: "All Truth Is Crooked, Time Itself Is A Circle" The Dwarf Represents The Spirit Of Gravity, And He Embodies The Herd Morality, "The Belittling Virtue" Which Is The Title Of Another Chapter From Part III. The Dwarf Can

Endure The Eternal Return Without Great Difficulties Because He Has No Aspirations; Unlike Zarathustra He Does Not Wish To Climb The Mountains That Symbolize Elevation And Solitude. In Two Unpublished Notes, From The Summer And The Fall Of 1883, Nietzsche Writes: The Doctrine Is At First Favored By The Rabble, Before It Gets To The Superior Men. The Doctrine Of Recurrence Will First Smile To The Rabble, Which Is Cold And Without Any Strong Internal Need. It Is The Most Ordinary Of Life Instincts, Which Gives Its Agreement First. Hence, The Content Of The Doctrine Is The Same, But Whereas The Dwarf Can Endure It (Because He Interprets It According To The Pessimistic Tradition For Which "Nothing Is New Under The Sun"), Zarathustra, Who Is The "Advocate Of Life" Regards The Eternal Return As The Strongest Objection To Existence, And As The Rest Of The Dream Suggests, He Does Not Yet Succeed In Accepting It After The Vision At The Gate Of The Instant, The Chapter Is Brought To An End By The Enigma Of The Shepherd. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO .How does the memory affect the GD complex or for that matter the gratification deprivation of an individual or a collective identity? Non-Markovian local-in-time master equations give a relatively simple way to describe the dynamics of open quantum systems with memory effects. Despite their simple form, there are still many misunderstandings related to the physical applicability and interpretation of these equations. Here, we clarify these issues both in the cases of quantum and classical master equations. Further introduction of the concept of a classical non-Markov chain signified through negative jump rates in the chain configuration. **E-M Laine et al 2012 J. Phys. B: At. Mol. Opt. Phys. 45 154004 doi:10.1088/0953-4075/45/15/154004 Local-in-time master equations with memory effects: applicability and interpretation.** Supposing That There Were Indeed An "Energy Of Contraction" Constant In All Centers Of Force Of The Universe, It Remains To Be Explained Where Any Difference Would Ever Originate. It Would Be Necessary For The Whole To Dissolve Into An Infinite Number Of

Perfectly Identical Existential Rings And Spheres, And We Would Therefore Behold Innumerable And Perfectly Identical Worlds COEXISTING [Nietzsche Underlines This Word Twice] Alongside Each Other. Is It Necessary For Me To Admit This? Is It Necessary To Posit An Eternal Coexistence On Top Of The Eternal Succession Of Identical Worlds? The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. Science is usually portrayed as dehumanizing. Brave New World epitomizes this fear. "The more we understand the world, the more it seems completely pointless" (Steven Weinberg). Certainly science can seem chilling when conceived in the abstract as a metaphysical world-picture. We may seem to find ourselves living in a universe with all the human meaning stripped out: Participants in a soulless dance of molecules, or harmonics of pointlessly waggling superstrings and their braneworld cousins. Nature seems loveless and indifferent to our lives. What rights have we to be happy? **(Brave New World: Aldous Huxley).** Religion can afford to be dogmatic but science cannot. Look at this passage which counters both scientists and philosophers alike: However, Hartmann Objects That The Regressive Movement Postulated By Schopenhauer Is Possible Only In Thought: It Remains Nothing More Than An "Ideal Postulate" With No Real Object And Which "Does Not Teach Us Anything About The Real Process Of The World That Unfolds In A Movement Contrary To This Backwards Movement Of Thought" (Hartmann, Philosophie Des Unbewussten, Third Edition (1871): 772). Hartmann Affirms That If Unlike Schopenhauer One Admits The Reality Of Time And Of The World Process, One Must Also Admit That The Process Must Be Limited In The Past And Therefore That There Must Be An Absolute Beginning. In Hartmann's Mind, Failure To Do So Would Result In Positing The Contradictory Concept Of An Accomplished Infinity: "The Infinity That From The Point Of View Of Regressive Thinking, Remains An Ideal Postulate, Which No Reality May Correspond To, Must, For The World, Whose Process Is, On The Contrary, A Progressive Movement, Open Up To A Determinate Result; And Here The Contradiction Comes To Light" (Hartmann

(1871): 772). What Really "Comes To Light" In This Passage Is The Fact That Hartmann Does Not Provide A Demonstration But A Petitio Principii. Indeed, The Concept Of The World Process Analytically Contains The Concept Of A Beginning Of The World. In All Rigors, It Is Therefore Impossible To Demonstrate These Concepts With Reference To Each Other. Secondly, Hartmann's View That One Is Bound To Accept The Reality Of The World Process Even If One Rejects The Ideality Of Schopenhauer's Time Is Mistaken. Hartmann Believes That If Time Is Real There Must Be A World Process With Both An Absolute Beginning And An Absolute End. Without Any Justification, Hartmann Jumps From Schopenhauer's Negated Time To Oriented Time. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. In the last years several theoretical papers discussed if time can be an emergent property deriving from ) quantum correlations. Here, to provide an insight into how this phenomenon can occur, authors present an experiment that illustrates Page and Wootters' mechanism of "static" time, and Gambini et al. subsequent refinements. A static, entangled state between a clock system and the rest of the universe is perceived as evolving by internal observers that test the correlations between the two subsystems. They implement this mechanism using an entangled state of the polarization of two photons, one of which is used as a clock to gauge the evolution of the second: an "internal" observer that becomes correlated with the clock photon sees the other system evolve, while an "external" observer that only observes global properties of the two photons can prove it is static.Phys. Rev. A 89, 052122 (2014)DOI:10.1103/PhysRevA.89.052122 arXiv: 1310.4691 [quant-ph] **Time from quantum entanglement: an experimental illustration Ekaterina Moreva.** Time is an entanglement .When the new ideas of quantum mechanics spread through science like wildfire in the first half of the 20th century, one of the first things physicists did was to apply them to gravity and general relativity. The results were not pretty. It immediately became clear that these two foundations of modern physics were entirely

incompatible. When physicists attempted to meld the approaches, the resulting equations were bedeviled with infinities making it impossible to make sense of the results. Then in the mid-1960s, there was a breakthrough. The physicists John Wheeler and Bryce DeWitt successfully combined the previously incompatible ideas in a key result that has since become known as the Wheeler-DeWitt equation. This is important because it avoids the troublesome infinites—a huge advance. But it didn't take physicists long to realise that while the Wheeler-DeWitt equation solved one significant problem, it introduced another. The new problem was that time played no role in this equation. In effect, it says that nothing ever happens in the universe, a prediction that is clearly at odds with the observational evidence. This conundrum, which physicists call 'the problem of time', has proved to be thorn in flesh of modern physicists, who have tried to ignore it but with little success. Then in 1983, the theorists Don Page and William Wootters came up with a novel solution based on the quantum phenomenon of entanglement. This is the exotic property in which two quantum particles share the same existence, even though they are physically separated. Entanglement is a deep and powerful link and Page and Wootters showed how it can be used (e) to measure time. Their idea was that the way pair of entangled particles evolves is a kind of clock that can be used to measure change. But the results depend on how the observation is made. One way to do this is to compare the change in the entangled particles with an external clock that is entirely independent of the universe. This is equivalent to god-like observer outside the universe measuring the evolution of the particles using an external clock. In this case, Page and Wootters showed that the particles would appear entirely unchanging—that time would not exist in this scenario. But there is another way to do it that gives a different result. This is for an observer inside the universe to compare the evolution of the particles with the rest of the universe. In this case, the internal observer would see a change and this difference in the evolution of entangled particles compared with everything else is an important a measure of time. This is

an elegant and powerful idea. It suggests that time is an emergent phenomenon that comes about because of the nature of entanglement. And it exists only for observers inside the universe. **Time is an entanglement** (Wikipedia). Danish religious philosopher Soren Kierkegaard carried out a systematic critique of the pretensions of reason and an abstract rationalism which he believed that the modern age was nurturing. Condemning reflection as a "danger" that ensnares people in logical delays and machinations, Kierkegaard compared it to a prison. Reflection is for him a form of captivity, a bondage which "can only be broken by [passionate] religious inwardness" (1978: 81). **Reflection seduces individuals into thinking its possibilities are "much more magnificent than a paltry decision"** (1978: 82). It leads them to act **"on principle,"** to dwell on the deliberation of the context of their actions and the calculation of their worth or outcome. Kierkegaard argues that this drives away feeling, inspiration, and spontaneity, all of which are crucial for true inner being and a vital relation to God. For Kierkegaard, as Nietzsche would later agree, genuine inner being (and culture) is characterized by the tautness and tension of the soul which characterizes passionate existence. But the "coiled springs of life relationships ... lose their resilience" in reflection (1978: 78) and "everything becomes meaningless externality, devoid of [internal] character" (1978: 62). Kierkegaard thus contributes to the development of an irrationalist tradition that has echoes in some later postmodern thought. Kierkegaard might well have agreed with his contemporary Fyodor Dostoyevsky, who wrote: "An intelligent [reflective] man cannot seriously become anything ... excessive consciousness is a disease" (1974: 3, 5). In an age overtaken by rules and regulations, genuine action -- which Kierkegaard assumes to be subjective and spontaneous -- is frustrated at every turn. Complaining that we are too "sober and serious" (1978: 71) even at banquets, Kierkegaard bemoans the fact that even suicides are premeditated (1978: 68)! "That a person stands or falls on his actions is becoming obsolete; instead, everybody sits around and does a brilliant job of bungling through with the aid of some reflection

and also by declaring that they all know very well what has to be done" (1978: 73). Thus, it is passion, not reflection, that guarantees "a decent modesty between man and man [and] prevents crude aggressiveness" (1978: 62). "Take away the passion and the propriety also disappears" (1978: 64). The ambiguity in the word "passion" may cause some confusion here. To say that the age and its individuals are "passionless" is not to say there are no emotions whatsoever, but rather that there is no true spiritual inwardness and depth, no intensively motivated action and commitment. It suggests that passion exists only in a simulated, pseudo-form, "the rebirth of passion" through "talkativeness" (1978: 64). "Chattering" for Kierkegaard gets in the way of "essential speaking" and merely "reflects" inconsequential events (1978: 89-99). Hence, in "the present age," emotions -- which in fact are all too pronounced -- have been transformed into negative forces. Anticipating Nietzsche's genealogy of the "slave revolt" in morality, Kierkegaard claims that the "enthusiasm" of the prior age of Revolution, a "positively unifying principle," has become a vicious "envy," a "negatively unifying principle" (1978: 81), a leveling force in its own right insofar as those lacking in talent and resources want to tear down those who have them. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS By Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html.** In November 2011, the United States government argued before the US Supreme Court that it wants to continue utilizing GPS tracking of individuals without first seeking a warrant. In response, Justice Stephen Breyer questioned what this means for a democratic society by referencing Nineteen Eighty-Four. Justice Breyer asked, "If you win this case, then there is nothing to prevent the police or the government from monitoring 24 hours a day the public movement of every citizen of the United States. So if you win, you suddenly produce what sounds like 1984... "[65]. What is to be said in unmistakable terms is that when it comes to national security and innocent citizens there is question of compromise even it be an infringement and encroachment of so called

privacy which has been made a word to subterfuge the activities that are far said to be fair. Quantum theory has provoked intense discussions about its interpretation since its pioneer days. One of the few scientists who have been continuously engaged in this development from both physical and philosophical perspectives is Carl Friedrich von Weizsäcker. The questions he posed were and are inspiring for many, including the authors of this contribution. Weizsäcker developed Bohr's view of quantum theory as a theory of knowledge

## ABSOLUTE STABILITY

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_9$ ) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9 - 9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2 - 2y_9\frac{(E_9 + n_9^2 y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right) - 4y_9(E_9 - n_9)$$

And has the solution

$$\frac{X_9}{M_9^2} > \frac{E_9}{(E_9 - 9)}[y_9(E_9 + n_9^2 y_9) + [y_9^2 (E_9 + n_9^2 y_9)^2 + 4y_9(E_9 - 9)(E_9 - n_9)]^{9/2}]$$

Given $E_9 > n_9$this lower limit to $X_9 / M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9 / M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9 < n_9$ $(ie., 2A_9 M_9^2 < X_9 - Y_9$)then (A2.9)and (A2.8)show that $Q_9$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2} > 2E_9[y_9 E_9 - (y_9 E_9 (9 + y_9)(E_9 - n_9)]^{9/2}]$$

And

$$E_9 > n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2/2M_9 (X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$and $X_9$and $Y_9$by $W_9^2$ the dependence upon *M* disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9$ *and* separate (A3.1)into real and imaginary parts

$$Z_9^2 + (M_9 + V_9) Z_9 \sin(Z_9 T_9) - (U_9 M_9 + Y_9) + (M_9(U_9 + V_9) - X_9)$$

$$\mathrm{X}\cos(Z_9 T_9) - 2M_9 V_9 \cos^2(Z_9 T_9) + M_9 V_9 = 0$$

$$(M_9 + U_9) Z_9 - (M_9 + V_9) Z_9 \cos(Z_9 T_9) + (M_9(U_9 + V_9) - X)$$

$$\mathrm{X}\sin((Z_9 T_9) - M_9 V_9 \sin(Z_9 T_9) = 0$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 = 1 - \cos Z_9 T_9$and $A_9 = U_9 - V_9$we obtain

$$4 M_9^2 V_9^2 W_9^4 + 2 M_9 V_9 W_9^3[(M_9 - V_9)^2 + A_9(M_9 + V_9) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2[M_9 - V_9)^2(2 A_9 M_9 + Y_9 - X_9) + (2A_9 M_9 - X_9)^2 - A_9(M_9 - V_9)(2A_9 M_9 - X_9)$$

$$+ M_9 V_9 (A_9^2 + 4(X_9 + Y_9))]$$

$$+ W_9[A_9^2 + 2X_9)(2 A_9 M_9 - X_9) + 2A_9 Y_9(M_9 + V_9) - M_9 A_9^3] + A_9^2(X_9 M_9 - Y_9) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9(2, U_9 V_9) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting $\cos(Z_9 T_9) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MDDULE NUMBERED SEVEN**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability

criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

## Morphological And Epistemological And Syllogistic Models For Which Nunny's Model Could Be Applied And Has Been Applied:

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production**

**Charge Asymmetry, $\Delta Y \equiv Yt - Yt^{-}$, Reflects The Asymmetry In $Tt^{-}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles

ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic

quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation

between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to

nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of

Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first

impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of

another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease. Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation

of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces

always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the

thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of

thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity

and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI

Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs” Outcomes—

And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities

Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \quad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

## DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t-T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\big(R(t)\big)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and haimovici model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

**TIME DELAY**

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Limitations of Predator prey analysis in immense. When we realize that world itself is Debit-Credit world and sometime credit-Credit world (See an earlier paper entitled the same nomeneclature.

Opposite limit of absolute stability. The result is a specific The Doctrine of Vibration awareness of an object of the form: 'here this is such.' While individual appearances do not lose their separate identity {svarupabheda) when they rest on a common basis, even so the particular object which appears according to its own characteristics (svalak$ana) is an individual reality in its own right. It is a different kind of appearance characterised by its association with the appearance of the specific location and time in which it is made manifest. The form of our experience is thus 'I now see this here'. But when we perceive each particular constituent appearance separately, each assumes a separate fixed function. Abhinava cites the following colourful example to illustrate how the various combinations of appearances account for the variety of experience: Thus even though the appearance of the beloved may manifest externally, it is as if far away in the absence of another appearance, namely, that of 'embracing'. So when the [appearing of the beloved] is associated with another appearance [namely that of 'far away'] the power (arthakriya) it formerly had of giving pleasure appears as its contrary. The form our experience assumes depends, not only on the nature of the object perceived, but also on personal factors entirely peculiar to ourselves. This theory explains this in two ways. In one sense, the object remains the same, but one or other of its constituent appearances

comes to the fore according to the inclinations of the perceiver. From another point of view, we can say that the perceived object is different for each perceiver according to the difference in the prominent appearance manifest to him. Abhinava, citing as an example a golden jar, illustrates how the same object appears differently to different perceivers according to the use they wish to make of it and to their state of mind: When a person who is depressed and feels that there is nothing [of value for him in the world] sees the jar, he merely perceives the appearance 'exists' [in the form of the awareness that] 'it is\ He is not conscious of any other [of its constituent appearances] at all. An individual who desires to fetch water [perceives] the appearance 'jar'. The man who simply wants something that can be taken somewhere and then brought back [perceives] the appearance 'thing'. The man who desires money [perceives] the appearance 'gold'. The man who desires a pleasing object [perceives] the appearance 'brightness' while he who wants something solid sees the appearance 'hardness'. These 'atomic events' or appearances emerge from the pure subject's consciousness and combine together to form a total event at each moment. Integral Monism of Kashmiri Saivism Daily life (vyavahara) goes on by virtue of this ever renewed flux of appearances. They are connected together and work towards a single unified experience because they appear within the field of consciousness of the universal subject. The aggregate of appearances arises in the [supreme] subject as do [sprouts in] a rice field. Even though each sprout germinates from its own seed, they are perceived as a collective whole. Appearances rest in this way within the universal subject. 'External- ity' is itself another appearance; it arises from a distinction between appearances and the individual subject. So, although all manifestation always occurs within the subject, it appears to be external due to the power of Maya which separates the individual subject from his object. This split must occur for daily life to be possible. Only externally manifest appearances can perform their functions; when they are merged within the subject and at one with him, they cannot do so. Daily life proceeds on the basis of the operation and withdrawal of the conditions necessary for fruitful action to be possible. Appearance in this sense represents the actualisation of a potential hidden in consciousness made possible by virtue of its dynamic, Spanda

nature which is both the flow from inner to outer and back as well as the power that impels it. The emergence from, and submergence into, pure consciousness of each individual appearance is a particular pulsation (visesaspanda) of differentiated awareness. Together these individual pulsations constitute the universal pulse (samanyaspandd) of cosmic creation and destruction. Thus, every single thing in this way forms a part of the radiant vibration {sphuratta, sphurana) of the light of absolute consciousness. Light and Awareness: The Two Aspects of Consciousness Absolute consciousness understood as the unchanging ontological ground of all appearing is termed 'Prakasa'. As the creative awareness of its own Being, the absolute is called ' Vimarsa '. Prakasa and Vimarsa — the Divine Light of consciousness and the reflective awareness this Light has of its own nature — together constitute the all-embracing fullness (purnata) of consciousness. The Recognition (pratyabhijfia) school of Kashmiri Saivism develops this concept of the absolute which finds its fullest expression in Utpaladeva's Stanzas on the Recognition of God. Even though neither of these two key terms appear in the Stanzas on Vibration or the Aphorisms of Siva, they recur frequently in their commentaries. Thus, although the original formulation of the Doctrine of Vibration differs from the theology of Recognition in this respect, it was extended in the course of its development to accommodate this concept of the absolute as well. This was possible, and quite justified, insofar as the absolute understood in Pratyabhijfia terms does not, as we shall see, differ essentially from that of the Spanda school. We can, as Kashmiri Saivites themselves have done, explain one in terms of the other. The Doctrine of Vibration Prakasa: The Light of Consciousness Prakasa is the pure luminosity' (bhdna) or 'self-showing' that constitutes the essence and ultimate identity (atman) of phenomena. That things appear at all is due to the light of consciousness, and their appearing (avabhasana) is itself this Light which bestows on all things their evident, manifest nature. Established in the light of consciousness everything appears there according to its own specific nature (svabhava). Anything that supposedly does not rest in this Light is as unreal as a sky-flower. 3 Thus, according to Rajanaka Rama, unlike the light of the sun, or any other light, this Light not only makes all things apparent, it is also their ultimate source. 4 Full of its

divine vibration the Light makes all things manifest and withdraws them into itself. This supra- temporal activity characterises it most specifically; devoid of it, it would be no better than an inert physical phenomenon. At the same time, this light is the conjunction (slesa) or oneness (aikdtmya) of its countless manifest forms, 6 and the collective whole (Sampina'ana) of all the categories of existence. The universe is nothing but the shining of the Light within itself. It is the radiant vibration (sphuratta) of this Light, the state (avast hand) in which consciousness becomes manifest. Although the Light shines as all things at all times and hence also makes their diversity manifest, 9 penetrating each object individually as well as collectively, it is not totally 'merged' (magna) or identified with the object so as to suffer any division within itself. Our experience of any object is of the form: T see this': it is not itself an object, but the manifest form the object assumes as a luminous principle of experience. The Light is ever revealed and can never be obscured; objectivity can never cast a shadow on the light of consciousness. The Stanzas on Vibration declare: That in which all this creation is established and from whence it arises is nowhere obstructed because it is unconditioned by [its very] nature. This Light is the highest reality (Paramartha). lity with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**DELAY DUE TO CONFEDERATION DEBIT TYPE MATURATION CREDIT TYPE CONSOLIDATION MATURATION**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \quad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9} + \frac{N_9 dG_9}{dR_9} - \frac{dB_9}{dR_9} > \frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9} + \frac{dM_9}{dR_9}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened**

**eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production Charge Asymmetry, $\Delta Y \equiv Yt - Y\bar{t}$, Reflects The Asymmetry In $Tt^{-}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton

particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results

in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely

dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact**

**most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the**

**thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is

nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the

world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease.

Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition

and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the

second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular

chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon

be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining

Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been

Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing

Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

**ABSOLUTE STABILITY**

$$L_9 = \frac{N_9 G_9}{2M_9} \frac{d(9nF_9)}{dR_9} (\geq 0) \quad (9a)$$

$$I_9 = -\frac{N_9 G_9}{2M_9} \frac{d(9nM_9}{dR_9} (\geq 0) \quad (9b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9 > L_9 - I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_9 > L_9 - I_9$, when $P_9^2/4 > I$ (8a)

$P_9 > L_9 - I_9 + \frac{(I_9 - P_9^2}{(L_9 - I)}$ when $P_9^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to

reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

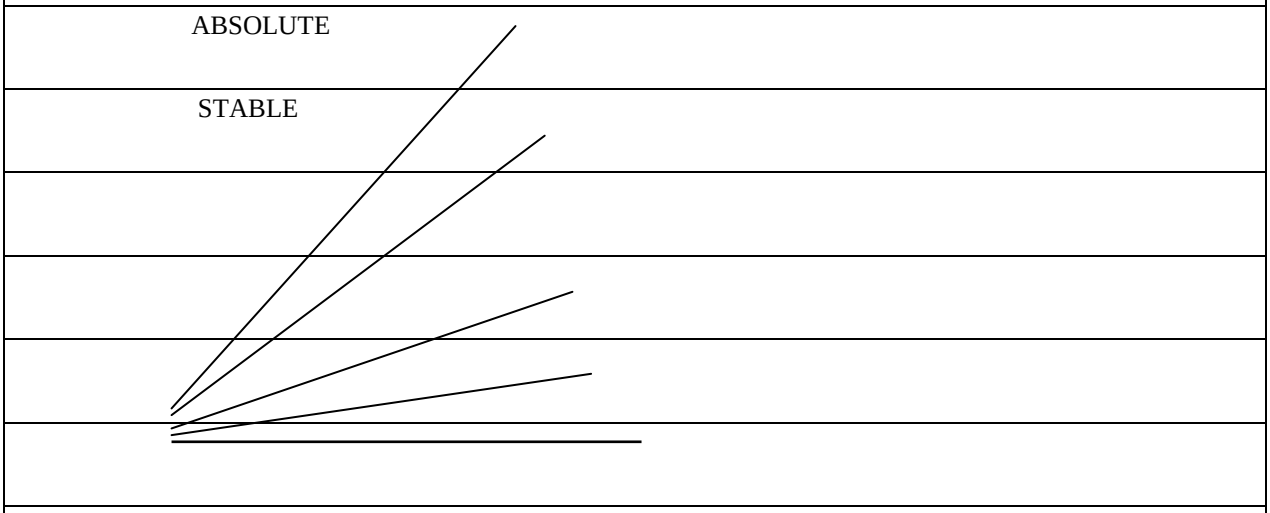



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori

reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

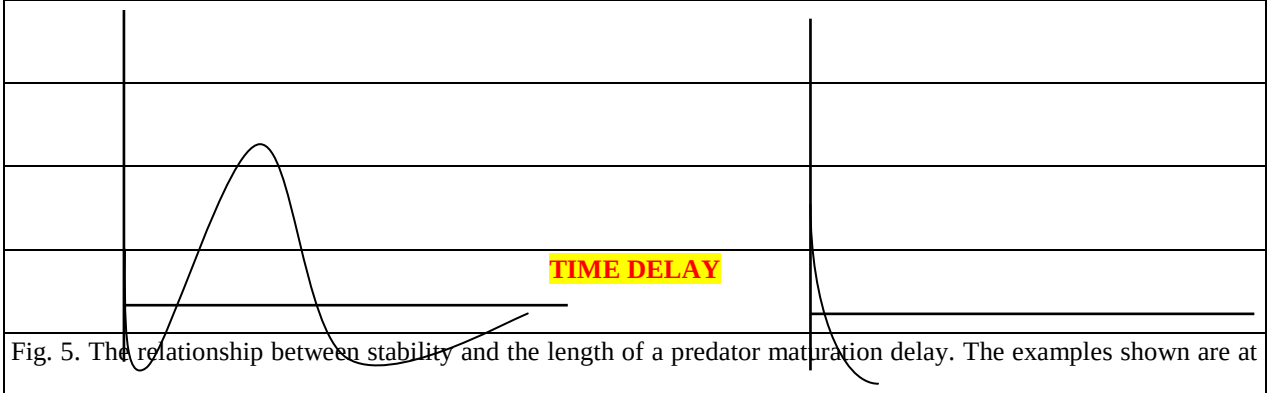



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

**DELAY IN BOTH PREDATOR AND PREY VIS A VIS DEBIT TYPES AND CREDIT TYPES**

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t - T_9)F_9\big(R_9(t - T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t - T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ and the effective prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting

$dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

Systems to which model has been applied also include: **Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI:

http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors.

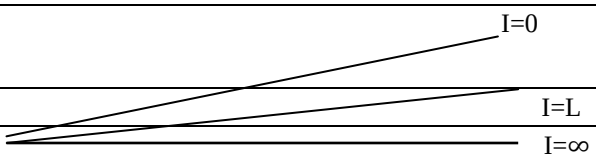
Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, ,**

**Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)**

TOTAL

ABSOLUTELY STABLE

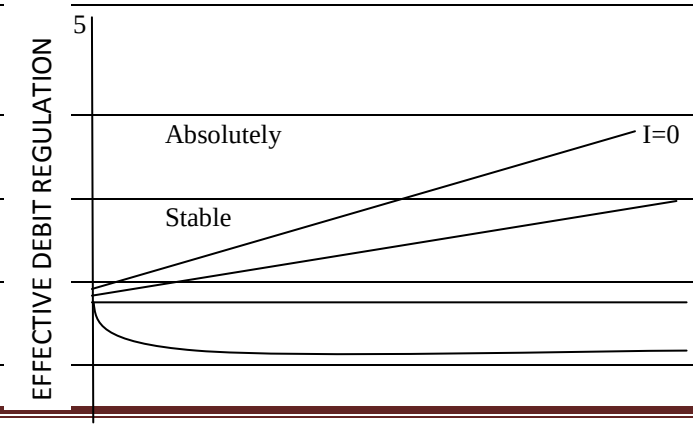



DELAYED DEBIT EFFECT

FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.



-50 10

**DELAYED CREDIT EFFECT (LACK OF BALANCE) CREDIT NON CUMSUMMATION:"**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

The following treatise bears ample testimony and infallible observatory to the fact that how Nunny's model can be applied to various cosmological, general relativity and economic problems.

## ONTOLOGICAL AND EPISTEMOLOGICAL FEATURES OF THE STUDY: OF NUNNY's MODEL

When they are known to be as they are, in the fullest sense of their existence, their phenomenal nature disappears leaving the ground of Being naked and accessible. This approach was validated by a critique of experience. The Vedanta established that space, time and the other primary categories of our daily experience can have no absolute existence. It was therefore necessary to make a distinction between relative truths — that accepted by the precritical common man — and an absolute truth discovered at a higher level of consciousness. The Saiva absolutist 17 rejects any theory that maintains that the universe is less than real. From his point of view a doctrine of two truths, one absolute and the other relative, endangers the very foundation of monism. The Kashmiri Saiva approach is integral: everything is given a place in the economy of the whole. It is equally wrong to say that reality is either one or diverse. Those who do so fail to grasp the true nature of things which is neither as well as both. "We do not" says Abhinavagupta, "base our contention that [reality] is one because of the contradictions inherent in saying that it is dual. It is your approach {paksa) that accepts this [method] While, if [duality and oneness] were in fact [to contradict each other], they would clearly be two [distinct realities]." The Vedantin, who maintains that non-duality is the true nature of the absolute by rejecting duality as only provisionally real, is ultimately landed in a dualism between the real and illusory by the foolishness of his own excessive sophistry (vacafadurvidya). Oneness is better understood as the coextensive unity (ekarasa) of both duality and unity. They are equally expressions of the

absolute. Gopinath Kaviraj says: According to Sarikara, Brahman is truth and Maya is inexplicable (anirvacaniya). Hence the [Advaitin's] endeavour to demonstrate the superiority of Advaita philosophy is turned against his own system. It tarnishes the picture of its philosophical perfection and profundity. He cannot accept Maya to be a reality, therefore his non-dualism is exclusive. The whole system is based on renunciation and elimination and thus is not all-embracing.... By accepting Maya to be Brahman (brahmamayi), eternal (nitya) and real (satyarupa), Brahman and Maya [in the Tantra] become one and coextensive. **The Doctrine of Vibration: An Analysis of the Doctrines and Practices of Kashmir Shaivism Mark S. G. Dyczkowski. Concept of being from Shaivite point of view:** Solutional behaviour and stability analysis of such systems is of paramount importance, notwithstanding the assumptions made which are not farfetched, in that it throws light on the fact whether they could be candidates for yet unsolved problems, providing Authetic ative determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication for the wide amplitudinal ramificatory usage of the thesis propounded. The world is now the other person's world, a foreign world that no longer comes from the self, but from the other. The other person is a "threat to the order and arrangement of your whole world...Your world is suddenly haunted by the Other's values, over which you have no control"(Satre: being and Nothingness) Sartre states that many relationships are created by people's attraction not to another person, but rather how that person makes them feel about themselves by how they look at them. This is a state of emotional alienation whereby a person avoids experiencing **their subjectivity by identifying themselves with "the look" of the other**. The consequence is **conflict**. In order to maintain the person's own being, the person must control the other, but must also control the **freedom of the other "as freedom**". These relationships are a profound manifestation of "bad faith" as for-itself is replaced with the other's freedom. The purpose of either participant is not to exist, but to maintain the other participant's looking at them. This system is often mistakenly called "love", but it is, in fact, nothing more than emotional alienation and denial of freedom. In

essence we are always simulating and selling ourselves at the cost of our freedom. Both as a script writer and participant consciousness, human beings experienceâ complete subjective feelings and sometimes absolute subjective feelings, which if remains unchecked might lead to pathological ramifications. (Italics mine) Perspicuous forbearance, Sophisticated seasoning, Partcipational observation, Background actuality, Existential worldliness, Existential strife, Predicational anteriorities, Character consonance, Ontological consonance, Primordial exactitude, Phenomenological correlates, Accolytish representation, Atrophied asseveration. Anamensial alienisms, Anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orinetatilities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state form the bastion, pillar, post, stylobate and sentinel of the classification scheme and doxa. appropriate and diagnostic, differentiating-discriminating, distinctive-distinguishing, emblematic- especial, essentially exclusive, fixed, idiosyncratic-inborn, inbred, indicative and individual, individualistic-inherent, individualizing and ingrained, inherent, innate, local, marked, native, normal, original, particularly- peculiar, personal, private, proper, regular, representative-singular, special, specific, symbolic, symptomatic- unique these form the classification and stratification measure that is suitable notwithstanding the constancy in many a case, as in gravity, which depends on the mass and distance for two given objects**.** Like Kant and Bergson, Deleuze considers traditional notions of space and time as unifying forms imposed by the subject. He therefore concludes that pure difference is non-spatio-temporal; it is an idea, what Deleuze calls "the virtual". (The coinage refers to Proust's definition of what is constant in both the past and the present: "real without being actual, ideal without being abstract.") Proust, Le Temps Retrouvé, ch. III: see the fourth line from the bottom of this page, or, in English translation, the thirteenth paragraph here: "I began to discover the cause by comparing those varying happy impressions which had the common quality of being felt simultaneously at the actual moment and at a distance in time, because of which common quality the noise of the spoon upon the plate, the unevenness of the paving-stones, the taste of the madeleine, imposed the past

upon the present and made me hesitate as to which time I was existing in. Of a truth, the being within me which sensed this impression, sensed what it had in common in former days and now, sensed its extra-temporal character, a being which only appeared when through the medium of the identity of present and past, it found itself in the only setting in which it could exist and enjoy the essence of things, that is, outside Time. [...] **Nothing but a moment of the past Much more perhaps; something which being common to the past and the present, is more essential than both. [...] a marvellous expedient of nature had caused a sensation to flash to me—sound of a spoon and of a hammer, uneven paving-stones**—simultaneously in the past which permitted my imagination to grasp it and in the present in which the shock to my senses caused by the noise had effected a contact between the dreams of the imagination and that of which they are habitually deprived, namely, the idea of existence—and thanks to that stratagem had permitted that being within me to secure, to isolate and to render static for the duration of a lightning flash that which it can never wholly grasp, a fraction of Time in its pure essence. When, with such a shudder of happiness, I heard the sound common, at once, to the spoon touching the plate, to the hammer striking the wheel, to the unevenness of the paving-stones in the courtyard of the Guermantes' mansion and the Baptistry of St. Mark's, it was because that being within me can only be nourished on the essence of things and finds in them alone its subsistence and its delight. It languishes in the observation by the senses of the present sterilised by the intelligence awaiting a future constructed by the will out of fragments of the past and the present from which it removes still more reality, keeping that only which serves the narrow human aim of utilitarian purposes. But let a sound, a scent already heard and breathed in the past be heard and breathed anew, simultaneously in the present and in the past, real without being actual, ideal without being abstract, then instantly the permanent and characteristic essence hidden in things is freed and our true being which has for long seemed dead but was not so in other ways awakes and revives, thanks to this celestial nourishment "While Deleuze's virtual ideas superficially resemble Plato's forms and Kant's ideas of pure reason, they are not originals or models, nor do they transcend possible experience; instead

they are the conditions of actual experience, the internal difference in itself. "The concept they [the conditions] form is identical to its object."[Desert Islands, p. 36.] A Deleuzean idea or concept of difference is therefore not a wraith-like abstraction of an experienced thing, it is a real system of differential relations that creates actual spaces, times, and sensations. [See "The Method of Dramatization" in Desert Islands, and "Actual and Virtual" in Dialogues. In fact it not as if the style and substance of science and philosophers have been entirely different. look at the conformality and consonance in the Gaussian Theory and Deleuze's statements on identity and the loss of it in the following passage…….That is, identity is a continuous process that finds its limit in exhaustion. As Deleuze will write, "…the logical relation of causality is inseparable from a physical process of signalling, without which it would not be translated into action. By 'signal' we mean a system with orders of disparate size, endowed with elements of dissymmetry; by 'sign' we mean what happens within such a system, what flashes across the intervals when a communication takes place between disparates. The sign is indeed an effect, but an effect with two aspects: in one of these it expresses, qua sign, the productive dissymmetry; in the other it tends to cancel it" (dr, 20). An actualized entity just is such a signal-sign system, remaining in perpetual communication with the virtual multiplicity it actualizes, constituting its own elements and identity. Such is Deleuze's account of what Badiou refers to as the operation of the "count-for-one"… an operation that is immanent to being itself and energetic in character. yet how are we to discern the local nature of ontological situations, the manner in which there is no global, overarching global situation in which all local situations are embedded as if parts of a whole? in part the above passage already responds to this question. if series must be brought into relation to communicate in order for causality to occur, and if the resonance of series is always a matter of chance, of a throw of the dice, like Lacan's objet a where no series enjoys primacy of model over copy, then all we have are local situations without a total situation. Unlike Heidegger's analysis of Dasein where Dasein is always being-in-the-world, Deleuze's simulacra or actualizations are always divergent, enjoying only local relations. Heidegger argues that in order to

understand the being of Dasein we must understand the manner in which Dasein is a being-in-the-world. By contrast, Deleuze's concept of multiplicity provides us with the principle of a new structuralism, of a new local ontology, that allows us to understand the immanent organization of a multiplicity without referring it to an embedding global space. as DeLanda so beautifully puts it in relation to the mathematician gauss, "…when gauss began to tap into these differential resources, a curved two-dimension surface was studied using the old Cartesian method: the surface was embedded in three-dimensional space complete with its own fixed set of axes; then, using those axes, coordinates would be assigned to every point of the surface; finally, the geometric links between points determining the form of the surface would be expressed as algebraic relations between the numbers. but gauss realized that the calculus, focusing as it does on infinitesimal points on the surface itself (that is, operating entirely with local information), allowed the study of the surface without any reference to a global embedding space. basically, gauss developed a method to implant the coordinate axes on the surface itself (that is, a method to implant the coordinate axes on the surface itself (that is, a method of 'coordinatizing' the surface) and, once points had been so translated into numbers, to use differential (not algebraic) equations to characterize their relations. as the mathematician and historian Morris Kline observes, by getting rid of the global embedding space and dealing with the surface through its own local properties, 'gauss advanced the totally new concept that a surface is a space in itself'" (intensive science and virtual philosophy, 11-12).according to the old Cartesian method, we can only outline the properties of a space by relativizing it to a global space in terms of which it is then mapped. by contrast, gauss is able to explore a space in terms of its intrinsic metric and organization as a local space, without referencing it to a whole of which it is conceptualized as a part. **The Image of Thought** "Image of thought "permeates both popular and philosophical discourse. According to this image, thinking naturally gravitates towards truth. Thought is divided easily into categories of **truth and error**. The model for thought comes from the educational institution, in which a master sets a problem and the pupil produces a solution which is either true or false. This image of the subject

supposes that there are different faculties, each of which ideally grasps the particular domain of reality to which it is most suited. In philosophy, this conception results in discourses predicated on the argument that "**Everybody knows..."** the truth of some basic idea. Descartes, for example, appeals to the idea that everyone can at least think and therefore exists. Deleuze points out that philosophy of this type attempts to eliminate all objective presuppositions while maintaining subjective ones. Deleuze maintains, with Artaud, that real thinking is one of the most difficult challenges there is. Thinking requires a confrontation with stupidity, the state of being formlessly human without engaging any real problems. One discovers that the real path to truth is through the production of sense: the creation of a texture for thought that relates it to its object. Sense is the membrane that relates thought to its other. Accordingly, learning is not the memorization of facts but the coordination of thought with a reality. "**As a result, 'learning' always takes place in and through the unconscious, thereby establishing the bond of a profound complicity between nature and mind".** Deleuze's alternate image of thought is based on difference, which creates a dynamism that traverses individual faculties and conceptions. This thought is fundamentally energetic and signifying: if it produces propositions, these are wholly secondary to its development. Philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; even tum tan tum;it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the vorticity of being merges with the positive use of the disjunctive synthesis, and

this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoin der for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for every hum of voices or a single silence for all the silences; a single vacuum for all the vacuousness consciousness should not be said without conjuring if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. The attributable to all states of bodies and states of affairs and the expressible of every proposition So univocity of consciousness means the identity of thé noematic attribute and that which is expressed linguistically and sense-fully. Univovity.We mean that it does not allow consciousness to be subsisting in aquasi state and but expresses in all pervading reality. (Reference: LOS: Deleuze) Deleuze sometimes spoke of getting up behind an author and "creating a monster", as if his reading methodology somehow consisted of a distortion of the philosopher's or author's thought. Certainly there is a case to be made for this with regard to the process of deterritorialization, where something is wrested from a territory and deterritorialized upon a new territory, like the animal paw that is deterritorialized from the earth and reterritorialize on the branch. however, it seems to me that the more interesting aspect of Deleuze's approach to other philosophers and art is not so much his "monstrous creations" (i seldom find them particularly monstrous), but rather their Gaussian or Riemannian style, where he explores them in terms of their own internal organization and metric, without reducing them to something alien such as history, society, biology (reductivism), or the signifier. What we find in Deleuze's approach to phenomena is a Gaussian technique. for instance, take Deleuze's books on cinema, his book on Francis bacon, or his study of sacher-masoch. In the first instance, Deleuze's carefully separates cinema from narrative and the signifier, studying it in terms of its specific organization pertaining to the production of images. in the case of bacon, Deleuze doesn't look for an underlying narrative or "meaning", but instead studies the manner in which bacon composes and

organizes his images and lines, both in terms of their production and actuality, so as to liberate a logic of sensation. finally, in approaching sacher-masoch, Deleuze allows Sacher-Masoch's novels to speak for themselves in terms of their desire and relation to pain, stalwartly refusing to reduce masochism to the complement of sadism. in each case, we have a local exploration of a "space" of multiplicities that is extremely precise. Sean Bowden the priority of events: **Deleuze's logic of sense sean Bowden, the priority of events: Deleuze's logic of sense, Edinburgh university press, 2011, 296pp., $40.00 (hbk), isbn 9780748643646 Reviewed by David Scott, Coppin state university and Gaussian spaces and multiplicities a note on Gaussian spaces and multiplicities posted by larval subjects under uncategorized.** What you "see" is not always what you "get." Many people mistakenly take their own visions literally without "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in the physical sense but in a psychological sense. The archetypes of the collective unconscious are arrayed behind the scenes of current worldwide conditions, of crisis and confusion. They mirror our own states back at us, whether we perceive them as such or interpret them plausibly or not. The noise of ordinary consciousness and beliefs drowns out the signal. Unconsciousness is the background of our ordinary awareness. Our organism is very much at the center of such effects. The organismic source is our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is (=) a meme-Scape. Pre-conceived concepts vie with (e&eb) structures, concepts with images. Like scientists who ignore (e) assumed truths, we leapfrog over (e&eb) our beliefs and personality deficits, claiming (eb) idiosyncratic imagination is (=) literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves (e) us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our

experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images. It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is (=) both. We inherit (e&eb) the structures that make our experience (=) what it is. But the structure itself is "empty," and each human culture "fills" it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about (e&eb) our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from

criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns,

nomenclature, and internal coherence of meaning and representation. n a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar some statements in Vedanta, it is not superfluous or redundant. It is evolutionary. Authetication determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication of the evolutionary process towards understanding is ever present in all, its thematic and discursive form. it is the relentless holding on to a theory this is despicable albeit the theory is not. Look at this paragraph. It is the 'I am the body/mind' belief that gives rise to the 'I am not (e) the world' belief. These two beliefs are (=) co-created."Consciousness projects (e&eb) the appearance of the mind, body and world by taking the shape of thinking, sensing and perceiving." "Attention is Consciousness with an object. When the object vanishes, attention simply remains what it always is, Consciousness.""There is (=) no purpose to (e) meditation. The purpose is already accomplished." "Everything that is experienced is experienced by, through, in and as Consciousness." "The seen cannot be separated from seeing and seeing cannot be separated from (e) Consciousness.""The Reality of any experience is not hidden in the appearance; it is expressed by the appearance." "Once we see that everything is Consciousness... Maya still dances, but it is a dance of love not seduction." There are some truly excellent sections

such as a long one taking, as a starting point, the following quotation from the painter Cézanne: "Everything vanishes, falls apart, doesn't it? Nature is always the same but nothing in her that appears to us lasts. Our art must render the thrill of her performance, along with her elements, the appearance of all her changes. It must give us a taste of her Eternity." Rupert begins: "That statement must be one of the clearest and most profound expressions of the nature and purpose of art in our era." And he goes on for some 14 pages to elaborate on this claim, examining the nature of the 'elements' Existence, appearance and Consciousness and their relationship

Some of the following concepts are modeled; some are not. For those which are not modelled kindly see one of the papers in the series. This note has the axiomatic predication to give a holistic view of the study. Here we talk of the characterstics of systems which satisfy the condition of cosmological constant or any other for that matter or variable or axiomatic predication. . There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit,

facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified

from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. Towards the end of circumvention of typing hundreds of superscripts and subscripts which would be a sardine squash and the concomitant operational difficulties, model is presented in piece meal of nine modules. Logarithms are to be taken in respect of those which incorporate more than one variable in bra-ket. Values of log (ab) and log (a+b) are readily available. Anti logarithm shall be taken at the end to predict, project, prognosticate the value of the variable. This is true for tensors, vectors and other variables too. Affirmational assertion, and explanation justification, statement of vindication, annotational commentary, and explicational glossary for each and every system changes and physical interpretation of results is one thing that is to be with earnest endeavour and feverish and febrile expediency. Editions never ever mean the same and identicality of parameters and this has been explicitly and unmistakably stated in the model apriori itself. Akin and analogous, cognate and concurrent, correspondent and congruent, comparable and complementary, synonymous and duple, tantamount and agnate, commensurate and correlative representation is only to highlight the importance and subterfuge the replication of work. It is my fervent solicitation to kindly bear with me for any lapses, notwithstanding the orchestrated efforts for a paper without any minor errors. Postulation predication, conclusive presumption, differential presuppositions, underscored decidedly axiomatic statement of the statement, equation form the bye word or the watch word in the aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement-argument, articulation, assertion, asseveration, choice of words, commentary and communication, declaration and definition, delivery-diction, elucidation-emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation tone and tenor of stratification. Now, how do find the reaction of systems to these singularities. You do the same thing a boss does for you. "Problematize" the events and see how you behave. I will resort to "pressure tactics". "intimidation of deriding report", or "cut in the increment" to make you undergo trials, travails and tribulations. I am happy to see if you improve your work; but may or may not be sad if you succumb to it and hang yourself! We do the same thing with systems. systems show conducive response, felicitous reciprocation or behave erratically with inner roil, eponymous radicality without and with blitzy conviction say like a solipsist nature of bellicose and blustering particles, or for that matter coruscation, trepidiational motion in fluid flows, or seemingly perfidious incendiaries in gormandizing fellow elementary particles, abnormal ebullitions, surcharges calumniations and unwarranted(you think so but the system idoes not!) unrighteous fulminations. Amount of emotion or affect that is appropriate or expected in an given form of interaction. Particular individuals and diffuse obligations (see c and d) are associated with affectivity, whereas contacts with many individuals (universalistic) in a bureaucracy may be devoid of emotion and characterized by affective neutrality. Affective neutrality may refer to self discipline and the deferment of gratification (eg Weber's spirit of capitalism). In contrast, affectivity may be associated with expressing emotions. Adams and Sydie also refer to affective neutrality being associated with ego control (p. 15). (Parsons: Wikipedia). So the point that is made here is "like we problematize the "events" to understand the human behaviour we have to "problematize" the events of systems to understand their behaviour. This statement is made in connection to the fact that there shall be creation or destruction of particles or complete obliteration of the system (blackhole evaporation) or obfuscation of results. Some systems are like "inside traders" they will not put signature atoll! How do you find they did it! Anyway, there are possibilities of a CIA finding out as they recently did! So we can do the same thing witn systems to. This is accentuation, corroboration, fortification, .fomentatory note to explain the various coefficients we have used in the model as also the dissipations called door. In the bank example we have clarified that various systems are individually conservative, and their conservativeness extends holisticallytoo.that one law is universal does not mean

there is complete adjudication of nonexistence of totality or global or holistic figure. Total always exists and "individual" systems always exists, if we donot bring Kant in to picture! For the time being let us not! Equations would become more ensorcelled and frenzied..... philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". In consciousness research, two rival sets of theories can be recognized: (A) Scientific material interpretations of consciousness are based on axioms that view consciousness in the context of highly advanced intentional processing of information in which subject-object relations evolve, and (B) humanistic interpretations of consciousness are based on axioms that view consciousness in the context of, say, "centered pulsations" that enable a conscious agent to act from his or her center of awareness. In this paper I will argue for the selection of axioms that favor humanistic interpretations of consciousness. The Pursuit of Autonomy Interdisciplinary Observations to Human Consciousness: Wautischer, Helmut (1993) the Pursuit of Autonomy. Interdisciplinary Observations to Human Consciousness: Social Neuroscience Bulletin, 6 (4). pp. 52-56.Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality; **CONSCIOUSNESS AND ITS UNIQUENESS June 12, 2012 at 8:40pm (See Deleuze Logic of sense, Wikipedia and Stanford encyclopedia for more details)** Atrocious contrivance and device, gratificational primogeniture, calamitous dodge and expedient, depraved, destructive, disastrous, execrable gambit and gimmick, great solace and succor, iniquitous, injurious, loathsome, low, maleficent, malevolent machination and maneuver, promethaleon of candor, frankness, honesty, honor, ingenuousness, innocence, openness, reality, sincerity, truthfulness, spiteful, stinking, ugly, unpleasant, unpropitious play and ploy, progenitor of jurisprudence and circumspection, racket and ruse, savvy, scam, stratagem, subterfuge, tactic and wile, forthrightness, honesty, truthfulness embodied and personified is the structural predisposition and dispensation of the complex. Nobody need have to take the Brahman-Anti Brahman agency to be a recalcitrant, repugnant, and refractory proposition. God gives and also takes away. That is the point made. ‘

And of course lesser gods, and God doth follow them and they follow God. That is why they are lesser gods. And wield the power on lesser mortals like us. When we mean gratification and deprivation we mean celestial Brahman Anti Brahman and Terrestial Brahman and Anti Brahman agencies producing such happiness or sadness by any means such as (blending- amalgamation, adulterant-adulteration, amalgam-amalgamation, blend- combination,

composite- compound, debasement-denaturant, fusion-hybrid, intermixture-reduction, amalgam- mixture, composite- compound, fusion-mishmash, aggregate- alloy, amalgamation-blend, commixture-composite, composition-compost, conglomerate-fusion, goulash-medley, mishmash-stew, blend-coalition, commixture-compound, federation-heating, integration-junction, melting-merger, synthesis-unification, assortment- combination, adulteration-alloy, assimilation-association, batter-blend, brewed combine, composite compound, concoction confection, conglomeration cross of money, orientation(dis), penury creation, murder, mayhem plunder, pillage, apocalypse, Armageddon in various forms and permutation and combination of Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness: with each action accounted and showing balance in the nature's general ledger as also actionlessness:e). All these are recorded in Nature's general ledger. Neuron DNA encapsulates all the actions passions, actions, interactions, transactions or the lack of it at various levels of individual, collective and cosmic levels. In celestial Brahman Anti Brahman ledger there is cessation of dualities. One more point to be made is if 5=3+2, then we can say 5 gormandizes 3 or 5 gobbles up 2. Quintessentially such a statement arises out of the difference of LHS and RHS term wise. It must be remembered despite all theoretical abstractions and generalizations be it proton mass or Higgs Boson, or topology of a space, it is notwithstanding the fundamentality of nature is just a number. Just everything in Physics and mathematics is just a number. Let us assume I buy rice in Bangalore (District) worth 5k.g. I can always divide 5k.g.s of rice in categories of Bangalore

(South), Bangalore (west), Bangalore (East) and Bangalore (South). On similar basis, there are billions of neutrons, protons and concomitant interactions which would satisfy the given numerical value. Example of gravity above makes things clear. Everything atleast is a function of time and someplace or some other coordinates, subordinates or superordinates. Mass of proton is 2.47 does not mean total mass of all protons is 2.47. It is this argument that is used in the classification scheme. By the very nature, self the unmonitoring but the witness consciousness agency identifies itself with something, be it a profession, an actor on the screen, when it does not have a identity. Such examples are legion, a Charlie Chaplin in "The Kid" or Raj kapoor in "Jagte Raho". These are people who just wanted to live and living without being put under surveillance every moment is the greatest part of their life. This is "Absolute Subjectivity" of Shiva. Such a state is always craving for identification for in pure consciousness all identities are lost. It is this "identitylessness" or having very minimal identity is what pervades most in India and this is due to expansion of individual consciousness. When that happens, be it education, or wealth, the "identity' becomes so well entrenched, they become the exact opposite of "pure consciousness"," impure consciousness" (Shakti). This is a state of "Absolute or relativistic objectivity". There was this case of a lady who went to such an extent of planning, she was sure that whatever she does must work out. When it did not, she lapsed in to "identity crisis". Such thing as this might happen in well developed countries where the competition is extremal and to be achieved at any cost. Realisation that one is not the doer but the witness consciousness is just like watching a movie in which one has acted. It is in this sense "advaita", "non duality' seems to be of cardinal and preeminent importance for the analytical treat ment of "identity' and "identity lessness", both happen in space time. "Multiverses" in this sense is also like a different film of an actor released at different theaters and it is in this sense we state that that every object in space time is a multiverse. You must have totally immersed in doing a "string theory" problem or watching a "film", oblivious of surroundings. Essentially creation (space time) also is what is happening in mind and honed in to you billion times. Identification of anything with a designation or signification is one that leads to

association within spacetime, which would only have pathological consequences and detrimental ramification as stated above or positive implications of "happiness" because of the identity itself. Search or 'show" has to go on! It gives a sense of consummation. a watching pot never boils; donot go to check nature; you are disturbing it; it will never show its true colours bhakti and Virakthi, gratification and deprivation, raga and dwesha is essential to live in space. It is the quantum interest thereof that produces energy to live .gratification and deprivation which form the bastion, pillar, post and stylobate and sentinel of all human actions depends on tams, rajas, and sattva of individual or insentient object. Verily tamas, rajas and satva are the characterstics of space itself, nay the behavioural pattern and attitudinal orientation. There are peaceful places like Vienna or Sringeri and there are violent places like others. Space is warped and woofed by the Akshara, the quantum information of primordial syllable Om. cosmic general ledger is the one vibration of which is brahman; quantum information possessions thereof form the sentinel of all transactions (Gargi to Yajnavalkya: my interpretation)adhocism, nonseqientiality of events, anti sequentiality of events, simultaneity of phenomenon, concurrent happenings , improbability of phenomenon happening itself are characterstics of brahman-anti brahman (Terrestial and celestial) are the cause of non predictability of happenings like in black hole, or in the areas of operation of brahman-anti brahman (Terrestial and celestial). brahman AntiBrahman actions, interactions, decisions, circumspection, adhocism, satisfies sensitive to initial conditions; topologically mixing; dense periodic orbits, the axiomatic predications and postulation alcovishness of chaos theory. Like gravity is a characteristic of space, so is tamas, rajas, sattva, with spaces knotted together by quantum information (for models see below) .accidents are preplanned actions of Terrestial and celestial Brahman AntiBrahman. All actions are perpetrated by them, be it train on fire or death. Belief systems play an important role in the determination of tamas, rajas, and satva state of human beings. There is neither cause nor effect. All actions are performed by Terrestial and celestial Brahman anti Brahman and it is off the record. psychic energy is negative equivalent to –mc^2, so that second law of thermodynamics is satisfied on a holistic basis vis-à-vis material energy

ParaBrahaman is infinite individual, collective and cosmic consciousnesses are infinite (concomitant general ledgers) for they are in the evolution mode with respect to each other in that order. Later emanates from the former as transactions emanate from the conservative state of assets and liabilities. Is not zero true for infinite conservative systems inside space and time? zero is the only consummation, consolidation and conservation. if fifty out of seventy members come out of the room at whatever time you come out, talk about the grand limbo of oblivion and hibernation of lifeless world you are going to and if you are not a celebrity, then it is not coincidence, it is coordination. Nature is playing the same trick on you. That a generalised law is there and you can measure does not make things less good albeit recalcitrant, refractory and rogue particles occur time and again, by and large. Entering in to all living beings through his constant part on the new moon night, he is born therefrom next morning. Absolute ParaBrahaman (zero) projects himself again and again on the screen of individual consciousness which appears of depairing despondent world. Some where the balancing act is done on karma. cosmic general ledger reflects upon itself to keep all transactional ties conservative (Brihadarnkayaka Upanishad: 1.5 14) for models see next paper through its divine power ( i interpret man creates god; it is attribution and ascription of various Shaktis would the formation of god takes place like space and time) the self assumes (eb) the form of a deity man can contemplate and venerate, even though siva, the pure subject, can never in fact be an object of meditation (adhyeya). Until we realise our true identity with Sankara, he is worshipped and conceived to be a reality alien to ourselves. While we are in the realm of creation, he too is a creation or mode or appearing of the absolute, manifest to us in meditation, through his freedom as an eternal, omniscient being. There is no gulf between the created and the uncreated creator: nothing in reality, although an object of knowledge, ceases to be Siva: this is the reason why meditation [on this or that aspect] of reality bestows its fruit. The world of the senses and mind appears to the well awakened (suprabuddha) as a theophany, an eternal revealing of god in his creation. The doctrine of vibration declares that "there is no state in word, meaning or thought, either at the beginning, middle or end, which is not Siva." 59 to utter any word is, in

reality, to intone a sacred formula.every act is a part of Siva's eternal cosmic liturgy, every movement of the body a ritual gesture (mudras), and every thought, god's thought. By what path are you not attainable? What words do not speak of you? in which meditation are you not an object of contemplation? What indeed are you not, o lord? 61 spandashakti, which accounts for the appearing of all things, is also the means by which deity in its many varied forms appears to man. Ksemaraja concludes: the ultimate object of worship of any theistic school differs not from the Spanda principle. The diversity of meditation is due solely to the absolute freedom of Spanda. 62 Sankara is not only the supreme object of devotion; as the static polarity of the absolute, he is the inner reality which holds together its Siva and Sakti outer manifestations. Phenomena are patterns of cognitions projected onto the surface of self-luminous siva-consciousness. There they become apparent, directly revealed to consciousness according to their manifest form. Siva is accordingly symbolized as the ground or surface of awareness, smooth and even like a screen (samabhittitalopama). 64 inscribed on this screen (kutfya) are the countless manifest forms which appear within it rendering it as diverse and beautiful as a fossil ammonite (falagrama). 65 Siva is the sacred ground upon which the cosmic mandalas are drawn, the absolute surface of inscription which bears the mark (cihna) of the universe. Abhinava writes: the variety of this world can only be manifest if the highest lord, who is essentially the pure light of consciousness, exists; just as a surface is necessary for a picture. if external objects were perceived in isolation then, because 'blue' and 'yellow', etc., are self-confined and the perceptions [we have of them] refer to their objects alone and so are insentient, mute and dumb in relation to one another . . . how would it be possible to be aware that an object is variegated? but just as depths and elevations can be represented by lines on a smooth wall, and we perceive [a female figure and think], similarly it is possible to be aware of differences in the variegated (contents of experience) only if all the diverse perceptions are connected together on the one wall of the universal light of consciousness. 66 Siva is the perfect artist who, without need of canvas or brush, paints the world pictures. The instant he imagines it, it appears spontaneously, perfect in every respect. The colours he uses are the varying shades

and gradations of his own Spanda energy and the medium his own consciousness. Intractable disaffect, Mutinous obdurate, Rebellious recalcitrant, Seditious refractory, Insubordinate insurgent, antagonistically antipodean, discrepant incompatible, Contumacious coadjutant, Conflicting converse, Unyielding obdurate, the antagonistic and anachronistic dispositional ties of Anti Brahman are in fact the sine qua non of Brahman only (See Vishwaroopa of Mahavishnu, which contains the polarised positives on RHS and negatives on LHS). The universe is coloured with the dye of its own nature (svabhava) by the power of Siva's consciousness (a//). Rajanaka rama says: homage to him who paints the picture of the three worlds, thereby displaying in full evidence his amazing genius {pratibha)\ to Sambhu who is beautiful with the hundreds of appearances laid out by the brush of his own unique, subtle and pure energy. 68 analogously, at the microcosmic level, all the cognitions and emotions, etc., which make up the individual personality form the outward flow of essentially introverted consciousness. They are specific pulsations (vise$aspanda) or aspects of the universal pulsation (samanyaspanda) of pure t consciousness. at the lower level, within the domain of maya, they represent the play in the fettered soul of the three primary qualities (gunas) or feeling-tones' which permeate to varying degrees his daily experience. These are: 1) sattva — the quality of goodness and luminosity which accompanies blissful experience both aesthetic and spiritual. 2) Rajas — the passion or agitation which oscillates between the extremes of 'light' and 'darkness' and characterises inherently painful experiences. 3) Tamas — the torpor and delusion which accompany states of inertia and ignorance. 69 the liberated soul recognises that these three are the natural and uncreated powers of pure consciousness. for him they are manifest respectively as: 1) Sankara 's power of knowledge (jnana) — the light of consciousness (prakasa); 2) the power of action (kriya) — the reflective awareness of consciousness (vimarsa)\ 3) the power of maya — which does not mean here the world of diversity, but the initial subtle distinction which appears between subject and object in pure consciousness. 70(doctrine of vibration: mark S.G. Dyczkowski) .to "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new

concepts and meanings that could not be included in the old system. at the heart of western metaphysics, for example, derrida finds the opposition between "speech" and "writing." this binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior in derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. The postmodern turn in philosophy: theoretical provocations and normative deficits by Steven best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html. You cannot see beyond consciousness for that is the limit of your storage of information or the basis for understanding further information. Turiya is empty bucket; it is from which prajna the measure of  arises; turiya measures the number of quantum information stores in other three states of waking, dreams, and dreamless deep sleep (Manduka Upanishad; my interpretation for models next paper).let there be path where angels tread: it is a force never born never dies the brahman an anti brahman verily there are parallel universes in (eb) in the ambit of space and time; the society of sinners and the sinned with the grammar of evil; parallel universes contain(e) multiverses in them; all physics, philosophy,  break down(e) in these universes; astronomical gratification ,deprivation, glorification, mortification, projections, axes, dimensions, rotations, foldings, concurrences, motions, shocks, are the axiomatic predications of multiverses and parallel universes; death is when self (the witness consciousness; register of individual general ledger) leaves the body; ostensibly parallel seem to have (e) an objective; parallel universes donot (e) have objectives; their paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common

attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, affection- appreciation, awareness- discernment, emotion-feeling, gut reaction, heart insight, intuition, judgment, keenness, perceptiveness, rationale, sensation, sense, sensitiveness, sensitivity, sentiment, susceptibility, taste, vibes, admiration, aesthetic sense, affection, appraisal, assessment, attraction, awareness, cognizance, commendation, comprehension, enjoyment, esteem, estimation, grasp, high regard, knowledge, liking, love, perceptional realization, recognitional regard, relish, respect, responsiveness, sensibility, sensitiveness, sensitivity, sympathy, understanding, valuation, acknowledgment, airing, baring, , confessional defenselessness, denudation- denunciation, disclosure, display, divulgence-exhibition, hazard, introduction, jeopardy, laying open, liability, manifestation, nakedness, openness, peril, presentation, publicity, revelation, risk, showing, susceptibility , susceptiveness,susceptivity, unfolding, unmasking, unveiling, vulnerability, vulnerableness, action, affection, appreciation and ardor, behavior, capacity, compassion, concern, cultivation and culture, delicacy, discernment and discrimination, emotion, empathy, faculty, fervor, fondness, heat, imagination, impression and intelligence, intensity and intuition, judgment, keenness- palpability, passion, pathos, pity, reaction, refinement, sensibility, sensitivity, sentiment, sentimentality , sharpness, spirit, sympathy, tangibility, taste, tenderness, understanding, warmth, fond memories, hearts and flowers, homesickness, longing, pining, reminiscence, ,schmaltz, sentimentality wistfulness, yearning satiation is only ostensible; their job is to get the job done; their business is to mind other's business ;all laws of physics and finance breakdown; it is antilaw,piscatorial Piratish and bubonic bucaneerishness; there is free transactions between in multiverses; when reflected on the screen of consciousness they appear as extant existential universe. Transactions are astronomical and gratification

deprivation, glorification and mortification incalculable; they have no barriers of region, religion, language which they create for others. theirs objectives are: corruption criminality, debasement debauchery, degradation evil immoral lewdness licentiousness, perversion sinfulness, vice viciousness, vileness wickedness, improbity corruption, atrocity decadence, degeneration degradation, depravity evil, immorality impurity, infamy iniquity, looseness lubricity, perversion profligacy, sinfulness turpitude,vice viciousness, vulgarity wickedness,disgrace bad reputation, abasement abuse, baseness black eye, blemish blur, brand comedown,contempt contumely, corruption culpability, debasement debasing, defamation degradation, derision disbarment, discredit disesteem, disfavor dishonor, disrepute disrespect, humbling humiliationignominy ill repute, infamy ingloriousness, meanness obloquy, odium opprobrium, pollution ,put-down reproach, scandal scorn, slander slight, slur spot, stain stigma, taint tarnish, turpitude venality, enormity Horribleness , abomination atrociousness, atrocity crime and the positive aspectionalities thereof .let there not be any incursions where evil mercenaries in the garb of mendicants thrive with adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatious gainsaying or conformingly pugnacious Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, execrable and goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness . statement for

Mephistophelean mercenaries philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality. Negative energy stabilising wormholes are considered candidates for (er) darkmatter. An analysis of negative energy might be a useful ingredient of the thoughts of travel at speeds more than the velocity of light. Peirce explains that an "imperfect continuum" possesses "topical singularities, or places of lower dimensionality" which break continuity, while "the parts of a perfect continuum have the same dimensionality

as the whole." (CP 4.642) Circular continua without singularities are time and space. (CP 6.210-212) By using math as his lens, Peirce envisioned singularities or lower dimensional holes in the imperfect cosmologic continuum, while maintaining that a probabilistic law of chance exists that entropy decreases **(Wikipedia).** On the side of philosophy**,** quite explicitly, Deleuze provides a critique of Husserl's "sleight-of-hand" change from a static categorical phenomenology to one supposedly more dynamically constitutive of reality. This a clear cut **chagrined circumvention**, **contraventional curbing**, **defeated disgruntlement**, **dissatisfactional downer**, **dragged failure**, **fizzled foil**, **grievance hindrance**, **impedimental irritation**, letdown, nonfulfillment, nonsuccess, obstruction, old one-two, **resent mental setback**, unfulfillment of the earlier theories makes Sartre contend that human existence is a conundrum whereby each of us exists, for as long as we live, within an overall condition of nothingness (nothing-ness)—that ultimately allows for free consciousness. But simultaneously, within our being (in the physical world), we are constrained to make continuous, conscious choices. It is this dichotomy that causes anguish, because choice (subjectivity) represents a limit on freedom within an otherwise unbridled range of thoughts. Subsequently, humans seek to flee our anguish through action-oriented constructs such as escapes, visualizations, or visions (such as dreams) designed to lead us toward some meaningful end, such as necessity, destiny, determinism (God), etc. Thus, in living our lives, we often become unconscious actors—Bourgeois, Feminist, Worker, Party Member, Frenchman, Canadian or American—each doing as we must to fulfill our chosen characters' destinies. However, Sartre contends our conscious choices (leading to often unconscious actions) run counter to our intellectual freedom. Yet we are bound to the conditioned and physical world—in which some form of action is always required. This leads to failed dreams of completion, as Sartre described them, because inevitably we are unable to bridge the void between the purity and spontaneity of thought and all-too constraining action; between the being and the nothingness that inherently coincide in our self. Deleuze insists on the absolute non-resemblance between what conditions (metastable, pre-individual plane of singularities) and what is conditioned

(actualized or the individuated 'thing'). Only then do "the conditions of the true genesis become apparent," Deleuze writes (Logic of Sense 105). Without addressing this issue, Deleuze's borrowings from Simondon, Lautman and structuralism remain untethered. Because of Bowden's decision to completely sidestep Deleuze's overt negotiations with phenomenology, his otherwise sophisticated commentary suggests a pedagogical narrowness. If, the omission of phenomenology strengthens Bowden's focus on the ontological priority of the event over substantial ontologies via the genealogy of figures and movements Deleuze outlines, it also incompletely articulates the problem of the sense-event as Deleuze formulates it. We have lost any idea of whom Deleuze takes to be his primary adversary, against whom he stakes out his position. This worry is most clearly manifested in Bowden's description of Deleuze's ontological prioritizing of events as offering a "transcendental ontology." For Bowden Deleuze's transcendental ontology defines events that "are ontologically prior to substances, 'all the way down'" (82). But Bowden provides a second definition: it is that the priority of events over worldly individuals means that the world of events "is not something external to the conditions of knowledge" (69). **SEAN BOWDEN the Priority of Events: Deleuze's Logic of Sense Sean Bowden, the Priority of Events: Deleuze's Logic of Sense, Edinburgh University Press, 2011, 296pp., $40.00 (hbk), ISBN 9780748643646. Reviewed by David Scott, Coppin State University Weltschmerz, agony, apprehension, blues, depression, dread, mid-life crisis, misgiving, nervousness, uneasiness, angst were expressed with such effulgent words…… aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement, argument, articulational assertion, asseverational choice of words, commentary and communication, declaration and definition, delivery, diction, elucidation, emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation about the human predicament is brought out in Being and Nothingness. Being and Nothingness offers a critique of Sigmund Freud's theory of the unconscious, based on the claim that consciousness is essentially self-conscious. Sartre also argues that**

**Freud's theory of repression is internally flawed. According to Sartre, in his clinical work, Freud encountered patients who seemed to embody a particular kind of paradox—they appeared to both know and not know the same thing. In response, Freud postulated the existence of the unconscious, which contains the "truth" of the traumas underlying the patients' behavior. This "truth" is actively repressed, which is made evident by the patients' resistance to its revelation during analysis. Yet what does the resisting if the patients are unaware of what they are repressing? Sartre finds the answer in what Freud calls the "censor". "The only level on which we can locate the refusal of the subject," Sartre writes, "is that of the censor.** The Holographic Principle, as conceived in current physics, applies to fields, and perhaps, even to more elementary entities, called strings and branes. These more elementary entities remain hypothetical at the time of this writing. There are many layers between the level of fields and that of neurons. We know that neurons are surrounded by a surface phospholipid membrane, which supports an electrochemical process that is fundamentally necessary to human experience. As the information supported by the membrane surface according to the Holographic Principle is proportional to its area, we should expect to find that the neural correlates of consciousness most clearly associated with those parts of the neuron that have the have a high proportion of neuronal surface area, and, in particular, receptive surface area. Dendrites account for an average of about 90% of the neuron is receptive surface area (Wong, 2002), and so would be expected to be the most important neural structures in information processes. We find this to be true **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science (Wikipedia).** Discordant-discrepant, dissonant-hard-line, implacable, incompatible, incongruous- inconsistent, inexorable-inflexible, inharmonious-intransigent, opposed, reluctant, unappeasable, uncompromising, unfriendly some of these expatiations and enumerations might sound, a more generalised interpretations is given by a Vedantin. Look at this arresting passage: The teacher continued to be silent. When addressed a second and third time he said: "I am teaching, but you do not follow. The Self is silence." The undetermined and unthinkable character of the Brahman is

a consequence of the absolute's eternal and immutable nature. To concede the existence of a real universe is, from the Vedantin's point of view, to posit the existence of a reality apart from the Brahman. Nor can we simply identify a real universe with the absolute unless we are prepared to compromise its unchanging, absolute status. The criterion of authenticity is immutability. Reality never changes; only that which is less than real can appear to do so. Reality is constant in the midst of change. What this means essentially is that there is change although nothing changes. This impossible situation is reflected in the ultimate impossibility of change itself. That which does not exist prior to its changing and at the end, after it has changed, must be equally non-existent between these two moments. Although the world of change appears to be real, it cannot be so. Change, according to the Vedantin, presupposes a loss of identity. Reality cannot suffer transformation; if it were to do so, it would become something else and the real would be deprived of its reality. The immortal can never become mortal, nor can the mortal become immortal. The ultimate nature of anything cannot change. Change of any sort is merely apparent (vivarta); the world of change and becoming is a false super- imposition (adhyaropa, adhyasa) on the absolute. In cosmic terms, the mistake (bhranti) consists of the supposition that the real Brahman is the unreal universe and the unreal universe is the real Brahman. In microcosmic terms, it is the mistake of falsely conceiving the body, mind or even one's personality to be the Self. In the same way as the image of a snake is falsely superimposed on a rope, similarly the universe is falsely projected onto the real substratum, the Brahman. Ignorance is not merely a personal lack of knowledge, but a cosmic principle. As such it is called "Maya," the indefinable factor (anirvacaniya) that brings this mistake in identity about. The reality status of this cosmic illusion is also indefinable: on the one hand it is not Brahman, the sole reality; on the other hand it is not absolutely non-existent like a hare's horn or the son of a barren woman. Brahman is the source of world appearances only in the sense of being their unconditioned ground or essential nature. The universe is false not because it has no nature of its own but because it does have one. Just as the illusion of a snake disappears when one sees that it is nothing but a rope, similarly cancellation (badha) of the empirically real occurs

when the absolute reality of the Brahman is realised. Thus, according to Vedanta, appearance implies the real, while the real need not imply appearance. To appear is essentially to appear in place of the real, but to be real is not necessarily to appear. All things -exist because the absolute exists. It is their Being. Thus the very existence of phenomena implies their non-existence as independent realities.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^\infty \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9 \left( U_9 - V_9 \int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9 \right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9 \int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$then given the constraints of $f(T_9), (A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 -$$

$$V_9 M_9 -]e^{-T_9\lambda_9} = 0$$

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9 Z_9 \sin(Z_9 T_9) - (A_9 M_9 + Y_9) + (A_9 M_9 + X_9)\cos(Z_9 T_9) = 0$$

$$(A_9 M_9) Z_9 - M_9 Z_9 \cos(Z_9 T_9) + (A_9 M_9 - X_9)\sin(Z_9 T_9) = 0$$

Quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High

Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) **Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP** This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) **Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP** In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference: Journal of Physics: Conference Series 70 **Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows:** System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol

Where $A_9 = U_9 - V_9$Using (A2.3), substitute for $\boldsymbol{Z}_9^2$ and $Z_9$in (A2.2):

$W_9^2 (X_9^2 - 3 M_9 A_9 X_9 - M_9^2 (X_9 - Y_9) + 2 M_9^3 A_9 + 2 A_9^2 M_9^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_9$=1-cos ($Z_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (X_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ X_9 - Y_9/2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \leq W_9 \leq 2$ *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (X_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\,M_9)^2\,(2A_9 M_9 - (X_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9 M_9^2 - 3M_9 X_9\right) + W_9\left(3M_9 A_9^2 - 2A_9 X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9+Y_9\right)$$

$$=2M_9^2W_9^2(M_9+\mathrm{A}_9)+M_9A_9^2W_9+\left(2A_9M_9-X_9+Y_9\right)(A_9W_9+M_9W_9^2$$

$$+(2-W_9)((2M_9X_9+M_9Y_9)W_9+A_9(X_9+Y_9)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range 0 $W_9$ 2.Defining

$Q_9\left(W_9,A_9\right)=a_9W_9^2+c_9$ *where* $a_9b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

To Heidegger writings we shall compose the model. Following is not still written model able form. For one of the models please see any one of the papers under the appellation: Ïn Excited Reveries

In his later writings Heidegger drops the term 'phenomenon' in preference for the verbal form 'phainesthai' in order to emphasize even more the actuality or presentational property of Being. Explaining this new form of the term he writes: "Being disclosed itself to the ancient Greeks as 'physis'. The etymological roots 'phy-' and 'pha-' designate the same thing: 'phyein', the rising-up or upsurge which resides within itself as 'phainesthai', lighting-up, self-showing, coming-out, appearing-forth." Heidegger contrasted his notion of phenomenon with semblance (Schein) and with appearing (Erscheinung). In the case of semblance a thing can show itself as that which it is not, as when fool's gold shows itself to be gold. The ancients always allied semblance with non-being. Heidegger points out, however, that semblances are grounded in

showings, and so does Abhinava. Both Heidegger and Abhinava consequently maintain that all semblances have a real basis and are to be treated as instances of phenomena along with the so-called real showing or manifestation of non-deceptive objects. So Heidegger states that: 'how- ever much seeming, just that much being'. Thus self-showing or appearing defines Being as phenomenon, but this definition of Being is as yet incomplete. Being is not only self-showing but 'logos' which Heidegger explains means 'discourse' (Rede) in the sense of 'apophansis': 'letting-be-seen'. Phenomenology, which according to Heidegger is the only correct study of Being, means 'letting-be-seen-that-which- shows-itself. This is true of Saiva Paramadvaita as well. The reality of the world demands recognition; we are forced to accept the direct presentation of the fact of our daily experience. As Abhinava says: "if practical life, which is useful to all persons at all times, places and conditions were not real, then there would be nothing left which could be said to be real" A thousand proofs could not make 'blue' other than the colour blue. The reality of whatever appears in consciousness cannot be denied. Objects appear; they do not cease to do so by a mere emphatic denial. The manifestation of an entity in its own specific form is a fact at one level of consciousness; it is real. The appearing of the same entity in the same form but recognised to be a direct representation of the absolute is also a fact, but at another level of consciousness. It is no more or less real than the first. 'As is the state of consciousness, so is the experience,' says Abhinava. Although the nature of the absolute is discovered at a higher level of consciousness, The Doctrine of Vibration nonetheless it presents itself to us directly in the specific form in which we perceive things; otherwise there would be no way in which we could penetrate from the level of appearing to that of its source and basis. Abhinava writes: Real is the entity iyastu) that appears in the moment of direct perception (sak$atkara), that is to say, within our experience of it. Once its own specific form has been clearly determined one should, with effort, induce it to penetrate into its pure conscious nature. All things are known to be just as they present themselves. The concrete actuality of

being known (pramiti), irrespective of content, is itself the vibrant (spanda) actuality of the absolute. Liberating knowledge is gained not by going beyond appearances but by attending closely to them. "The secret," Mahesvarananda says, "is that liberation while alive (jivanmukti) is the profound contemplation of Maya's nature." No ontological distinction can be drawn between the absolute and its manifestations because both are an appearing (dbhasa), the latter of diversity and the former of 'the true light of consciousness which is beyond Maya and is the category Siva'. Those who have attained the category of Pure Knowledge above Maya and have thus gone beyond the category of Maya, see the entire universe as the light of consciousness . . . Just as the markings [on a feather] are nothing apart from the feather, the feather [is nothing apart from] them, similarly, when the light of consciousness is manifest, the whole group of phenomena is manifest as the light of consciousness itself. Within the sphere of Maya, every entity's 'own nature' (svabhdva) corresponds to its specific manifest form. Accordingly it is defined as that which distinguishes it from all else and from which it never deviates. Above the sphere of Maya, that is, above the level of objectivity, is the domain of the subject. At this level, everything is realised to be part of the fullness of the experience r and hence no longer bound by the conditions which impinge on the object. Here the part is discovered to be the whole, that is, consciousness in toto. In this sphere beyond relative distinctions, the yogi realises that (all) the categories of existence are present in every single category. m The yogi experiences every individual particular as the sum total of everything else. He recognises that all things have one nature and that every particular is all things. This is the 'essence' (sard) or co-extensive unity (samarasya) of all things. Integral Monism of Kashmiri SaivismWe have established that reality is manifest according to how [and the degree in which] the freedom of consciousness reveals it and that [this freedom] is the womb of all forms. Just as 'sweetness' is present in its entirety in every atom of the sugarcane, so each and every atom [of the universe] bears within itself the emanation of all things. This is

the level of consciousness in which the absolute reflects on itself realising to its eternal delight and astonishment (camatkara) its own integral nature. The reality of the world of diversity is not denied, but experienced in a new mode of awareness free of time and space in the eternal omnipresence of the Here and Now. [Phenomenal forms of awareness] such as 'this [exists]', born of the colouring [imparted to the absolute] by the limitations engendered by the diversifying power of time {kalakalana) also emanate within the Supreme Principle. There [at that level], Fullness {purriata) is the one nature [of all things] and so everything is omnipresent; otherwise, associated with division (khancjlana), the Fullness [of the absolute] would not be full. The content of absolute consciousness consists of diverse appearings (abhdsa) which, because they are manifest through it in this way, do not compromise the wholeness of consciousness. Everything we perceive is a momentary collocation of a number of such manifestations which combine together like 4 a row of altar lamps' (dipavalT) to form the single radiant picture of the universe. The individual objects which constitute the universe are specific collocations of such 'atomic' appearings. Together they form a single unified particular which appears according to its own defining features (svalaksand). A jar, for example, consists of a number of appearances such as 'round', 4 fat\ 'earthen', 'red', etc., which together discharge a single function (arthakriya), in this case, that of carrying the appearance 'water'. They unite with each other much as the scattered rays of a lamp come together when focused, or as the various currents of the sea together give rise to waves. Atomic appearings can combine in any number of ways, provided that they are not contrary to one another as established by the dictates of natural law (niyati). An appearance of 'form', for example, cannot combine with that of 'air'. Insofar as they share a common basis (samanyadhikaranya), a given cluster of appearances appears as a single whole. This common basis is the most prominent member of the group; the appearance 'jar' is such in the example quoted above. Any one appearance in a cluster may assume a more important or subordinate role.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_9 \square - b_9 \square 0$$

And

$$b_9^2 \square 4a_9 c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9 in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\ X_9$ ,

$Let\ Y_9 = \ y_9\ X_9$And try the substitution $A_9 = X_9/(2 + A_9\ y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + \ n_9 y_9)$ we have

$$W_9^2\ y_9\left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^2\left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9\,X_9\left[\frac{X_9\left(1+n_9y_9\right)}{M_9^2\left(2+n_9y_9\right)^2}+2\,y_9\left(n_9-1\right)+\frac{X_9^2\left(9+y_9\right)}{M_9^2\left(2+n_9y_9\right)}\right]=0.\ (A2.12)$$

Since $n_9$ and $A_9$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9=\left(2+n_9y_9\right),$ *the condition*$(A2.99?)$*can be written as*

$$\left[\frac{X_9\left(E_9-9\right)}{E_9^2}\right]^2\left[\frac{X_9^2}{M_9^2}-4\,y_9\,E_9^2\frac{X_9}{M_9^2}+4y_9\left(n_9-2\right)E_9^3\right]$$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9/W_9^2)$ can be expressed as

$$2\,E_9\left[y_9E_9+\left(y_9\,E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}\right]>\frac{X_9}{M_9^2}>2E_9\left[y_9E_9-y_9E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}]$$

dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique

for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) **Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP** Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) **Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP** Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a

cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) **Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP** Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service

placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) **Dynamic Service Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP** Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by $H*(X,(\wedge*TX)\otimes g)$, where X is the target hyperk\"ahler manifold. We further give a mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) **A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van**

**Raamsdonk1 and Nathan Seiberg2** study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. **Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP** Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy **The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP** Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the

present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. **ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP** Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition. Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb) the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537

© 2016 American Physical Society **Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP** Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant ($\Delta\alpha/\alpha$) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and $\Delta\alpha/\alpha$ to (e&eb) this conclusion. However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts $z\lesssim 3$, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the $\Lambda$ cold dark matter model value ($j\Lambda CDM=1$); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in $\alpha(z)$, linked to (e&eb) a linear coupling ($\zeta$) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order $|\zeta|\sim 10-4$, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in $\Delta\alpha/\alpha$ (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society **The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP** Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation

rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void

Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems, authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society **Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP** Authors measure the star formation properties of (e) two large samples of

galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the **star formation** rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a **function of environment**. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]**.** Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size $\sqrt{n} \times \sqrt{n}$? The problem is that $\sqrt{n}$ time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time $O(\sqrt{n})$ for $d \geq 3$, or $O(\sqrt{n}\log 5/2n)$ for $d = 2$. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search

tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ∗This work was mostly done while the author was a **PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship and by CIAR**. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 **Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP** Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can

lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 **Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP** Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. **Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP** Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e)

sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) **Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP** The relationship between the foundations of mathematics and (e&eb) physics is a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb) flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling

field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) **Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP** The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar

and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph] (or arXiv: 1503.05600v1 [math-ph] for this version) **Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP** This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs

mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) **Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP** Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and

derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to **the fine structure constant**. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) **Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP** Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor goes to 0 as the (e&eb) time goes to 0, the

big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph] (or arXiv: 1306.4613v1 [quant-ph] for this version) **Effects of gauge theory based number scaling on geometry Paul Benioff STOP** In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x,

outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph]

(or arXiv: 1211.3381v2 [quant-ph] for this version) **Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP** Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value,

\int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances, the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) **Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP** The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include

(e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) **Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP** Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) **New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP** This work is based on (e) the field of reference frames based on (e) quantum

representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects:Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc) Journal reference:Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) **A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul B**

**ABSOLUTE STABILITY**

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_9$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9-9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2-2y_9\frac{(E_9+n_9^2y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right)-4y_9(E_9-n_9)$$ (A2.13)

And has the solution

$$\frac{X_9}{M_9^2}>\frac{E_9}{(E_9-9)}[y_9(E_9+n_9^2y_9)+[y_9^2(E_9+n_9^2y_9)^2+4y_9(E_9-9)(E_9-n_9)]^{9/2}]$$

Given $E_9>n_9$this lower limit to $X_9/M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9/M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9<n_9$ $(ie.,2A_9M_9^2<X_9-Y_9)$then (A2.9)and (A2.8)show that $Q_9$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2}>2E_9[y_9E_9-(y_9E_9(9+y_9)(E_9-n_9)]^{9/2}]$$

And

$$E_9>n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2/2M_9(X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$ and $X_9$ and $Y_9$ by $W_9^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

**Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability

of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e) long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007 STOP. David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer P** examine the possibility of coherent reversible information transfer between solid-state superconducting qubits and (e&eb) ensembles of ultracold atoms. Strong coupling between these systems is mediated by (e&eb) a microwave transmission line resonator that interacts near resonantly with (e&eb) the atoms via (e&eb) their optically excited Rydberg states. Solid-state qubits can then be used to (e) implement rapid quantum logic gates, while (e) collective metastable states of (e) the atoms can be employed for (e) long-term storage and (e&eb) optical readout of quantum information. **Reversible state transfer between superconducting qubits and (e&eb) atomic ensembles David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer Phys. Rev. A 79, 040304(R) – Published 14 April 2009 STOP.** Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures have been proposed. An

attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. **Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier** demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in (eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. **Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP.** In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown

to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , **Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally** frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. **Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg** propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states

into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. **Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011 http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan** proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 **Qubit transfer between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller** propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. **Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys.**

**Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich** report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. **Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi** investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. **Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev** investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer.

Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. **L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP.** In recent years **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P have** investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P** apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys. Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. **Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S.**

**Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP.** Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert **Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe** realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.**220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2** [quant-ph] for this version) **Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP.** Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. **Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker** describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue

from over (e&eb) a thousand plant and animal samples. **NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP.** Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this behavior, enhancing and accelerating (eb, eb+) the growth of correlations. **Michael Kastner** gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. **Michael Kastner** tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) **Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP.** A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. **Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E. Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91,**

**075435 – Published 26 February 2015 STOP.** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still

problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to

(e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP** Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that

(eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035

**Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime**

**C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013**

**STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and

cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP** Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex

Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) **A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen**

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9 and$ separate (A3.1)into real and imaginary parts

$$Z_9^2 + \left(M_9 + V_9\right) Z_9 \sin\left(Z_9 T_9\right) - \left(U_9 M_9 + Y_9\right) + (M_9\left(U_9 + V_9\right) - X_9)$$

$$\mathrm{X}\cos\left(Z_9 T_9\right) - 2M_9 V_9 cos^2\left(Z_9 T_9\right) + M_9 V_9 = 0$$

$$(M_9 + U_9)\, Z_9 - (M_9 + V_9)\, Z_9\ cos\left(Z_9 T_9\right) + \left(M_9\left(U_9 + V_9\right) - X\right)$$

$$\mathrm{X}\sin\left(\left(Z_9 T_9\right) - M_9 V_9 \sin\left(Z_9 T_9\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 =$ $1 - cos\, Z_9\ T_9$and $A_9 = U_9 - V_9$we obtain

$$4\, M_9^2 V_9^2 W_9^4 + 2\, M_9 V_9\ W_9^3[(M_9 - V_9)^2 + A_9\left(M_9 + V_9\right) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2\left[M_9 - V_9\right)^2\left(2\, A_9\ M_9 + Y_9 - X_9\right) + \left(2A_9\ M_9 - X_9\right)^2 - A_9\left(M_9 - V_9\right)\left(2A_9\ M_9 - X_9\right)$$

$$+ M_9 V_9\,(A_9^2 + 4\left(X_9 + Y_9\right))]$$

$$+ W_9\left[A_9^2 + 2X_9\right)\left(2\, A_9\ M_9 - X_9\right) + 2A_9 Y_9\,(M_9 + V_9) - M_9 A_9^3] + A_9^2\left(X_9\ M_9 - Y_9\right) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9\left(2, U_9 V_9\right) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting cos $(Z_9 T_9) = -1\ in\ (A3.3)$

We have applied the model to Kashmiri Shaivism. Whose elaborate rituals is enough to make a indivual raise from individual consciousness to cosmic consciousness which we have called Cosmic Consciousness.

It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) **Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP** The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic

Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** The group of unimodular transformations on the roots of polynomials over GF(2) is considered, and those polynomials with (e&eb) symmetries in (eb) the unimodular group are identified. The cross-correlation function between (e&eb) two maximum-length linear shift register sequences of (e) the same degree is shown to be (=) computable as an explicit linear transformation, in (eb) matrix form, on (e&eb) either one of the sequences, regarded as (=) a vector. The underlying vector space is (=) the "cyclotomic algebra," generated by (e) the cyclotomic cosets, or by what Gauss termed the **"periods" of the cyclotomic equation**. A variety of numerical examples are worked in detail. **Information Sciences Volume 1, Issue 1, December 1968, Pages 87-109 Theory of transformation groups of polynomials over GF(2) with applications to linear shift register sequences Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more link. Solomon W. Golomb** stop Authors study the P−V criticality and phase transition in (eb) the extended phase space of charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of (e) the system and its conjugate quantity is (=) the thermodynamic volume of the black hole. The black holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb) no P−V criticality and phase transition appear, while (e) for the black holes with a spherical horizon, even when the charge of the black hole is absent, t(eb) he P−V criticality and the small black hole/large black hole phase transition will appear, but it happens only in d=5 dimensions; when the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to have (e) the P−V criticality

and (e&eb)the small black hole/large black hole phase transition, there exists (ebb) an upper bound for the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the **Gauss-Bonnet coefficient** and Q is (=) the charge of the black hole. Authors calculate the critical exponents at the critical point and find that (eb) for all cases, (eb) they are the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09(2013)005 Cite as: arXiv:1306.6233 [gr-qc] (or arXiv:1306.6233v4 [gr-qc] for this version) **P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in IR but with (e&eb) improved UV behaviors. Here, authors study an IR modification which breaks ``softly'' (e) the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Authors obtain the black hole and cosmological solutions for **``arbitrary'' cosmological constant** that represent (eb) the analogs of the standard Schwarzschild-(A)dS solutions which can be asymptotically (A)dS as well as flat and authors discuss their thermodynamical properties. They also obtain solutions for (e) FRW metric with (e&eb) an arbitrary cosmological constant. Authors study its implication to (e&eb) the dark energy and find that it seems to be consistent with (e&eb) current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP.** The Einstein equations with (e&eb) a negative cosmological constant admit (e) black hole solutions which are asymptotic to (e) anti-de Sitter space. Like black holes in asymptotically flat space, these solutions have (e) thermodynamic properties including (e) a characteristic temperature and intrinsic entropy equal to (=) one quarter of the area of the event horizon in Planck units. There are however some important differences from (e) the asymptotically flat case. A black hole in anti-de Sitter space has (e) a minimum temperature which occurs when (e) its size is of the order of the characteristic radius of (e) the anti-de Sitter space. For larger black holes the red-shifted

temperature measured at (eb) infinity is greater. This means that such black holes have (e) positive specific heat and can be in stable equilibrium with (e&eb) thermal radiation at a fixed temperature. It also implies (eb) that the canonical ensemble exists for asymptotically anti-de Sitter space, unlike the case for asymptotically flat space. One can also consider the microcanonical ensemble. One can avoid (e) the problem that arises in asymptotically flat space of (e) having to put the system in (eb) a box with unphysical perfectly reflecting walls because (e) the gravitational potential of anti-de Sitter space acts as a box of finite volume. Communicated by A. Jaffe **Communications in Mathematical Physics December 1983, Volume 87, Issue 4, pp 577-588 Thermodynamics of black holes in anti-de Sitter space S. W. Hawking, Don N. Page STOP** In previous work authors introduced the concept of black-hole entropy, which they identified with (e&eb) the surface area of the black hole in question expressed in units of (e) the Planck length squared. Authors suggested that the appropriate generalization of the second law for a region containing a black hole is (=) that the black-hole entropy plus (+) the common entropy in the black-hole exterior never decreases. Here authors establish the validity of this law for the infall of an entropy-bearing system into (e&eb) a much larger and more massive generic stationary black hole. To do this authors determine (eb) a general lower bound for the increase in black-hole entropy, and an upper bound for the entropy of the system, while allowing for quantum effects at each stage. In passing authors show that the generalized second law is a statistical law which becomes (=) over-whelmingly probable in the limit of a macroscopic system. Authors also consider briefly more general situations. Finally, we give two simple examples of predictions made by the generalized second law for (e) black-hole formation processes. Received 17 September 1973 DOI:http://dx.doi.org/10.1103/PhysRevD.9.3292 **Generalized second law of thermodynamics in black-hole physics Jacob D. Bekenstein Phys. Rev. D 9, 3292 – Published 15 June 1974 STOP** The Lovelock gravity consists of (e) the dimensionally extended Euler densities. The geometry and horizon structure of (e) black hole solutions could be quite complicated in (eb) this gravity; however, authors find that some thermodynamic quantities

of the black holes like the mass, Hawking temperature and entropy have (e) simple forms expressed in terms of (e) horizon radius. The case with black hole horizon being a Ricci flat hypersurface is (=) particularly simple. In that case the black holes are (=) always thermodynamically stable with (e&eb) a positive heat capacity and their entropy still obeys (e&eb) the area formula, which is no longer valid for (e) black holes with positive or negative constant curvature horizon hypersurface. In addition, for black holes in (eb) the gravity theory of Ricci scalar plus a 2n-dimensional Euler density with (e&eb) a positive coefficient, thermodynamically stable small black holes always exist in (eb) D=2n+1 dimensions, which are absent in (eb) the case without (e) the Euler density term, while (e) the thermodynamic properties of the black hole solutions with (e&eb) the Euler density term are qualitatively similar to those of black holes without (e) the Euler density term as D>2n+1. Copyright © 2004 Published by Elsevier B.V. **Physics Letters B Volume 582, Issues 3–4, 4 March 2004, Pages 237–242 A note on thermodynamics of black holes in Lovelock gravity Rong-Gen CAI STOP** There are a number of similarities between black-hole physics and (e, eb, =) thermodynamics. Most striking is the similarity in the behaviors of black-hole area and (e&eb) of entropy: Both quantities tend to increase (eb+) irreversibly. In this paper we make this similarity the basis of (e) a thermodynamic approach to black-hole physics. After a brief review of the elements of the theory of information, authors discuss black-hole physics from the point of view of information theory. Authors show that it is natural to introduce (eb) the concept of black-hole entropy as (=) the measure of information about (e&eb) a black-hole interior which is (=) inaccessible to an exterior observer. Considerations of simplicity and consistency, and dimensional arguments indicate (eb) that the black-hole entropy is equal to (=) the ratio of the black-hole area to the square of the Planck length times a dimensionless constant of (e) order unity. A different approach making use of the specific properties of Kerr black holes and of concepts from information theory leads to (eb) the same conclusion, and suggests (eb) a definite value for the constant. The physical content of (e) the concept of black-hole entropy derives from (e) the following generalized version of the second law: When common entropy goes down a

black hole, (eb) the common entropy in (eb) the black-hole exterior plus the black-hole entropy never decreases. The validity of this version of the second law is supported by (e) an argument from information theory as well as by several examples. Received 2 November 1972 DOI:http://dx.doi.org/10.1103/PhysRevD.7.2333 **Black Holes and Entropy Jacob D. Bekenstein Phys. Rev. D 7, 2333 – Published 15 April 1973 STOP** The effects of power-law Maxwell field on the Van der Waals like phase transition of higher-dimensional dilaton black holes are probed in detail. It is shown that the Smarr relation gains corrections due to (e) the effects of (e) both the power-law Maxwell field and (e&eb) the dilaton field while (e) thermodynamic volume is exactly the same as (=) that of Einstein-Maxwell-dilaton black holes. Authors successfully derive the analytic solutions of (e) critical point and carry out (e&eb) some check to ensure that these critical quantities are (=) positive. It is shown that the constraint on the parameters turns out to be (=) 0<\alpha^2<1, which is tighter than that in (eb) the non-extended phase space. It is also shown (eb) that these critical quantities and (e&eb) the ratio P_cv_c/T_c are affected by (e&eb) the power-law Maxwell field. Moreover, critical exponents are found to coincide with those of other AdS black holes, showing (eb) the powerful influence of mean field theory. Comments: references added Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1601.05500 [gr-qc] (or arXiv:1601.05500v2 [gr-qc] for this version **Effects of power-law Maxwell field on the Van der Waals like phase transition of higher dimensional dilaton black holes Jie-Xiong Mo, Gu-Qiang Li, Xiao-Bao Xu STOP** A consistent and unified picture for (e) critical phenomena of charged AdS black holes in f(R) gravity is drawn in this paper. Firstly, we investigate the phase transition in canonical ensemble. Authors derive the explicit solutions corresponding to (e&eb) the divergence of CQ. The two solutions merge into (e&eb) one when (e) the condition Qc=−13R0−−−√ is satisfied. The curve of specific heat for Q<Qc has two divergent points and can be divided into (e&eb) three regions. Both the large radius region and the small radius region are (=) thermodynamically stable with (e&eb) positive specific heat while (e) the medium radius region is unstable with (e&eb) negative specific heat. However, when Q>Qc, the specific heat is

always positive, implying (eb) the black holes are locally stable and no phase transition will take place. Secondly, both the T−r+ curve and T−S curve f(R) AdS black holes are investigated and they exhibit Van der Vaals like behavior as the P−v curve in the former research. Critical physical quantities are obtained and they are consistent with (=) those derived from the specific heat analysis. Authors carry out numerical check of Maxwell equal area law for the cases Q=0.2Qc, 0.4Qc, 0.6Qc, 0.8Qc. The relative errors are (=) amazingly small and can be negligible. So the Maxwell equal area law holds for (e) T−S curve of f(R) black holes. Thirdly, authors establish geometrothermodynamics for (e) f(R) AdS black hole to examine the phase structure. It is shown that the Legendre invariant scalar curvature R would diverge exactly where (e) the specific heat diverges. To summarize, the above three perspectives are consistent with (e&eb) each other, thus providing (eb) a unified picture which deepens the understanding of (e) critical phenomena of (e) f(R) AdS black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1602.01251 [gr-qc] (or arXiv:1602.01251v1 [gr-qc] for this version) **A consistent and unified picture for critical phenomena of f(R) AdS black holes Jie-Xiong Mo, Gu-Qiang Li, Yu-Cheng Wu STOP** The state equation of the charged AdS black hole is reviewed in (e&eb) the T−r plane. Thinking of the phase transition, the T−S, P−V, P−ν graphs are plotted and then the equal area law is used (e) in the three cases to get (eb) the phase transition point (P, T). The analytical phase transition point relations for (e) P-T of charged AdS black hole has been obtained successfully. By comparing the three results, authors find that the equal area law possibly cannot (e) be used directly for P−ν plane. According to the T−S, P−V results, they plot (e&eb) the P−T−Q graph and find that for a highly charged black hole a very low temperature condition is required for (e) the phase transition. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI:10.1140/epjc/s10052-015-3641-0 Cite as: arXiv: 1503.07658 [gr-qc] (or arXiv:1503.07658v1 [gr-qc] for this version) **A note on Maxwell's equal area law for black hole phase transitionShan-Quan Lan, Jie-Xiong Mo, Wen-Biao Liu STOP** Recently, extended phase space thermodynamics of Lovelock AdS black holes has been of great interest. To provide

insight from (e) a different perspective and gain (e) a unified phase transition picture, non-extended phase space thermodynamics of (e) (n+1)-dimensional charged topological Lovelock AdS black holes is (=) investigated detailedly in (eb) the grand canonical ensemble. Specifically, the specific heat at (eb) constant electric potential is calculated and phase transition in the grand canonical ensemble is discussed. To probe the impact of the various parameters, they utilize (e) the control variate method and solve (e) the phase transition condition equation numerically for the case k=1, −1. There are two critical points for (e) the case n=6, k=1 while there is only one for other cases. For k=0, there exists no phase transition point. To figure out the nature of phase transition in the grand canonical ensemble, authors carry out an analytic check of the analog form of (e) Ehrenfest equations proposed by Banerjee et al. It is shown that Lovelock AdS black holes in (eb) the grand canonical ensemble undergo (e&eb) a second order phase transition. To examine the phase structure in (eb) the grand canonical ensemble, authors utilize the thermodynamic geometry method and calculate both the Weinhold metric and Ruppeiner metric. It is shown that for both analytic and graphical results that the divergence structure of (e) the Ruppeiner scalar curvature coincides with (=) that of the specific heat. Research provides (eb) one more example that Ruppeiner metric serves as (=) a wonderful tool to (e) probe the phase structures of (e) black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C75 (2015) 5, 211 DOI:10.1140/epjc/s10052-015-3434-5 Cite as: arXiv: 1503.01956 [gr-qc] (or arXiv: 1503.01956v1 [gr-qc] for this version **Non-extended phase space thermodynamics of Lovelock AdS black holes in grand canonical ensemble Jie-Xiong Mo, Wen-Biao Liu STOP** To provide an analytic verification of the nature of phase transition at the critical point of P−V criticality, the original expressions of Ehrenfest equations have been introduced directly. By treating the cosmological constant and (e&eb) its conjugate quantity as (=) thermodynamic pressure and volume respectively, authors carry out analytical check of (e) classical Ehrenfest equations. To show that approach is universal, authors investigate not only higher-dimensional charged AdS black holes (e&eb), but also rotating AdS black holes. Not only are the examples

of Einstein gravity shown, but also the example of modified gravity is presented for (e) Gauss-Bonnet AdS black holes. The specific heat at constant pressure CP, the volume expansion coefficient α and the isothermal compressibility coefficient κT are found to diverge exactly at the critical point. It has been verified that both Ehrenfest equations hold at (eb) the critical point of P−V criticality in (eb) the extended phase spaces of (e) AdS black holes. So the nature of the critical point of P−V criticality of AdS black holes has been demonstrated analytically to be (=) a second-order phase transition. These results are consistent with (e&eb) the nature of liquid-gas phase transition at the critical point. In this sense, extant research would deepen the understanding of (e) the relations of AdS black holes and (e&eb) liquid-gas systems. Moreover, successful approaches to introduce (eb) the original expressions of Erhenfest equations directly into (e&eb) black hole phase transition research demonstrate again (eb) that black hole thermodynamics is closely related (e&eb) to classical thermodynamics, which allows (eb) us to borrow techniques from (e) classical thermodynamics to investigate the thermodynamics of black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 89, 084057 (2014) DOI: 10.1103/PhysRevD.89.084057 Cite as:arXiv: 1404.3872 [gr-qc] (or arXiv:1404.3872v1 [gr-qc] for this version) **Ehrenfest scheme for P−V criticality of higher dimensional charged black holes, rotating black holes and Gauss-Bonnet AdS black holes Jie-Xiong Mo, Wen-Biao Liu STOP** Aiming at a unified phase transition picture of (e) the charged topological black hole in (eb) Ho\v{r}ava-Lifshitz gravity, authors investigate this issue not only in (eb) canonical ensemble with (e&eb) the fixed charge case but also in grand-canonical ensemble with (e&eb) the fixed potential case. Authors firstly perform the standard analysis of the specific heat,(e&eb) the free energy and the Gibbs potential, and (e&eb) then study its geometrothermodynamics. It is shown that the local phase transition points (eb) not only witness the divergence of (e) the specific heat, but also witness the minimum temperature and the maximum free energy or Gibbs potential. They also witness the divergence of (e) the corresponding thermodynamic scalar curvature. No matter which ensemble is chosen, the metric constructed can

successfully produce (eb) the behavior of the thermodynamic interaction and (e&eb) phase transition structure while other metrics failed to (e) predict the phase transition point of (e) the charged topological black hole in former literature. In grand-canonical ensemble, authors have discovered the phase transition which has not been reported before. It is similar to (e) the canonical ensemble in which the phase transition only takes place when k=−1. But it also has its unique characteristics that the location of the phase transition point depends on (eb) the value of potential, which is different from (e) the canonical ensemble where (e) the phase transition point is (=) independent of the parameters. After an analytical check of Ehrenfest scheme, authors find that the new phase transition is (=) a second order one. It is also found that the thermodynamics of (e) the black hole in Horava-Lifshitz gravity is quite different from that in (eb) Einstein gravity. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP1310(2013)056 DOI:10.1007/JHEP10(2013)056 Cite as: arXiv:1404.2497 [gr-qc] (or arXiv:1404.2497v1 [gr-qc] for this version) **A unified phase transition picture of the charged topological black hole in Horava-Lifshitz gravity Jie-Xiong Mo, Xiao-Xiong Zeng, Gu-Qiang Li, Xin Jiang, Wen-Biao Liu STOP** To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case $\beta\to\infty$ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness condition of (e) entropy. It is shown

that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure authors have discovered. It is shown that Legendre invariant thermodynamic scalar

curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws, leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions, geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged AdS black holes by considering the case of the RN-AdS black holes

surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase transition. For the case $\omega_q=-2/3$ authors derive analytic expressions of (e) critical physical quantities, while (e) for cases $\omega_q\neq-2/3$ appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this

version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=) chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy

Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed

quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions, while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma?**

All this consists of the act of consciousness alone", says Vasubandhu, "because unreal entities appear, just as a man with defective vision sees unreal hair or a moon, etc."He points to dreams as examples of purely subjective constructs which appear to be objective realities. The apparent reality dreams possess is not derived from any concrete, objective world, but merely from the idea of objectivity. While the Yogacara does not say that an idea has, for example, spatial attributes, it does have a form manifesting them. While he agrees with the Saiva idealist that

appearances have no independent existence apart from their appearing to consciousness, he maintains that for this reason they are unreal. The creativity of consciousness consists in its diversification in many modes having apparent externality; it is not a creation of objects. While the Kashmiri Saivite agrees that the world is pure conscious- ness alone, he maintains that it is such because it is a real creation of consciousness. The effect is essentially identical with the cause and shares in its reality. Matter and the entire universe are absolutely real, as 'congealed' (sty ana) or 'contracted' (samkucita) forms of consciousness. "This God of consciousness", writes Ksemaraja, "generates the universe and its form is a condensation of His own essence (rasa)" m By boiling sugarcane juice it condenses to form treacle, brown sugar and candy which retains its sweetness. Similarly, consciousness abides unchanged even though it assumes the concrete material form of the five gross elements. The same reality thus abides equally in gross and subtle forms. Consequently no object is totally insentient. Even stones bear a trace (vasana) of consciousness, although it is not clearly apparent because it is not associated with the vital breath (prana) and other components of a psycho-physical organism. Somananda goes so far as to affirm that physical objects, far from being insentient, can only exist insofar as they are aware of themselves as existing. The jar performs its function because it knows itself to be its agent. Indeed, all things are pervaded by consciousness and at one with it and hence share in its omniscience. Thus, Siva, Who perceives Himself in the form of physical objects, is the one ultimate reality. "The jar knows because it is of my nature", writes Somananda, "and I know it because I am of the jar's nature. I know because I am of Sadasiva's nature and He knows because He is of my nature; Yajnadatta [knows] because he is of Siva's nature and Siva [knows] because He is of Yajnadatta's nature". Integral Monism of Kashmiri Saivism Everything in this sense is directly perceived by absolute conscious- ness, and this direct perception (pratyaksa) unifies the knowable into a single, undivided whole. This is the central concept behind a doctrine originally expounded by Narasimha called 'the non-dualism of direct perception' (pratyaksadvaita). This states that consciousness is essentially perceptive and that its perception of all things operates throughout the universe. Insofar

as phenomena are clearly evident (sphufa) to us, everything is directly perceived by absolute consciousness, with which our individual consciousness is identical. This direct perception unfolds everywhere; the one true reality, it is alone and without companion or rival (nihsapatna). Even though it remains one, it can, by its very nature, perceive distinctions (bheda) between one entity and another, without this engendering any division within it. We distinguish between two entities in empirical terms on the basis of their mutual exclusion (anyonyabhava). The relative distinction {bheda) between them is essentially the perceived difference between their respective characteristics. Despite this difference they are united within the purview of a single cognition insofar as they are equally both manifest appearances. This cognition is the undivided essence (rasa) or 'own nature' (svabhava) of both. Encompassed by the 'fire of consciousness', there is no essential difference between them. Just as when an emerald and ruby reflect each other's light, the ruby is reddish-green and the emerald greenish-red, similarly everything is connected with everything else as part of the single variegated (vicitra) cognition of absolute consciousness. Mahesvarananda writes: The Supreme Lord's unique state of emotivity (asadharanabhava) is the outpouring of pure Being (mahasatta). It is manifest as the brilliance (sphuratta) of the universe which, if we ponder deeply, [is realized to be] the single flavour (ekarasa) of the essence of Beauty which is the vibration of the bliss of one's own nature. In this way all things are in reality one although divided from the one another sharing as they do the 'single flavour' (ekarasa) of the pure vibration of consciousness. Kashmiri Saiva Realism Kashmiri Saivism as a whole has been variously called a form of 'realistic idealism', 'monistic idealism', ‘idealistic monism' and 'concrete monism'. It is easy to understand why Kashmiri Saivism is The Doctrine of Vibration said to be 'idealistic' and 'monistic', but in what sense is it also 'realistic*? The answer to this question is of no small importance in trying to under- stand the central idea behind its metaphysics and the fundamental importance of the concept of Spanda, in this seemingly impossible marriage between monistic idealism and pluralistic realism. The Kashmiri Saiva approach understands the world to be a symbol of the absolute, that is, as the manner in which it presents itself to us. Again we can

contrast this view with that of the Advaita Vedanta. The Advaita Vedanta understands the world to be an expression of the absolute insofar as it exists by virtue of the absolute's Being. Being is understood to be the real unity which underlies empirically manifest separateness and as such is never empirically manifest. It is only transcendentally actual as 'being-in-itself. The Kashmiri Saiva position represents, in a sense, a reversal of this point of view. The nature of the absolute, and also that of Being, is onceived as an eternal becoming (satatodita), a dynamic flux or Spanda, 'the agency of the act of being'. It is dentified with the concrete actuality of the fact of appearing, not passive unmanifest Being. Appearance (abhasa) alone is real.ppearing (prakasamanatva) is equivalent to the fact of being (astitva). Ksemaraja writes in his commentary on the Stanzas on Vibration: Indeed, all things are manifest because they are nothing but manifestation. The point being that nothing is manifest apart from manifestation The absolutely unmanifest, from this point of view, can have as little existence as the space in a lattice window of a sky-palace. Nay, even less, because even that space can appear as an imagined mage manifest within consciousness. Everything is real according to the manner in which it appears. Even an illusion is in this sense real, insofar as it appears and is known in the manner in which it appears. The empirical and the real are identical categories of thought. As Abhinava says: Thus this is the supreme doctrine (upanisatf), namely that, when- ever and in whatever form [an entity] appears, that then is its particular nature. Perhaps at this stage a brief comparison with Heidegger's ideas might prove to be enlightening and not altogether out of place. According to Heidegger's phenomenology of Being, reality is intelligible in a two-fold Integral Monism of Kashmiri Saivism manner as 'phenomenon' and 'logos'. Heidegger defines what he means by 'phenomenon' as: "that-which-shows-itself. The manifest . . . phenomena are then the collection of that which lies open in broad daylight or can be brought to the light of day — what the Greeks at times implicitly identified as 'ta onta' (the things-which-are)".

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb) heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna**

**Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version) **Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb) strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel

period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar (=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI: 10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills (SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear

differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black**

**holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-

relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb)

the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this

version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstate thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11 August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of

entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.** Characterizing the behaviour of strongly coupled

quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at $\lambda=\lambda c$, the correlation length diverges as $\xi\sim|\lambda-\lambda c|-\nu$, while (e) the gap between the ground state and first excited state vanishes as (e) $\Delta E\sim\xi-z\sim|\lambda-\lambda c|z\nu$. Exactly at the critical point $\lambda=\lambda c$, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as $\omega\sim|k-k0|z$. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-**

**equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004 STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element.

**Model has been applied for the following results:**

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| **Searches for Sterile Neutrinos with the IceCube Detector M. G. Aartsen et al. (IceCube Collaboration) Phys.**<br><br>**Rev. Lett 117, 071801 – Published 8 August 2016 Physics logo See Viewpoint: Hunting the Sterile Neutrino** |
| **(1) The IceCube neutrino telescope at the South Pole has measured the atmospheric muon neutrino spectrum as (=) a function of zenith angle and (e&eb) energy in the approximate 320 GeV to 20 TeV range, to search for (e) the oscillation signatures of (e) light sterile neutrinos.** |
| **(2) No evidence for anomalous νμ or ν¯μ disappearance is observed in (eb) either of two independently developed analyses, each using one year of (e) atmospheric neutrino data.** |
| **(3) New exclusion limits are placed on (e&eb) the parameter space of (e) the 3+1 model, in which muon antineutrinos experience (e) a strong Mikheyev-Smirnov-Wolfenstein-resonant oscillation.** |
| **(4) The exclusion limits extend to (e) sin22024≤0.02 at Δm2∼0.3 eV2 at the 90% confidence level.** |
| **(5) The allowed region from (e) global analysis of appearance experiments, including (e) LSND and MiniBooNE, is excluded at approximately the 99% confidence level for (e) the global best-fit value of \|Ue4\|2. Received 5 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.117.071801 © 2016 American Physical Society** |
| **Synopsis: Searching for Majorana Neutrinos August 16, 2016The KamLAND-Zen collaboration has run the most sensitive search to date for a radioactive decay that could reveal whether neutrinos are Majorana fermions. Synopsis figureKamLAND-Zen Collaboration** |
| **(6) Are neutrinos and (=) antineutrinos the same particles, as hypothesized by (e) physicist Ettore Majorana?** |
| **(7) If true, this "Majorana" nature of neutrinos could explain (eb) why neutrinos have (e) mass and why the Universe has (e) more matter than antimatter.** |
| **(8) Using an underground neutrino detector in Japan, the KamLAND-Zen collaboration has carried out the most sensitive search to date for a process called neutrinoless double-beta decay, which could confirm (eb) Majorana's hypothesis.** |
| **(9) The results set stringent constraints on the possible properties of (e) Majorana neutrinos.** |
| **(10) Double-beta decay is (=) a nuclear process in which two neutrons turn into (e&eb) two protons (or vice versa), emitting (eb) two antineutrinos.** |
| **(11) If neutrinos are Majorana fermions, (eb) the two antineutrinos could annihilate each other, and some double-beta decays would emit (eb) no antineutrinos.** |
| **(12) However, neutrinoless decay is believed to be (=) extremely rare. Searches of the process using (e) large amounts of the 35 naturally occurring (eb) isotopes that can undergo (e&eb) double-beta decay (such as xenon-136) have so far come up empty.** |

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| **(13) The KamLAND-Zen collaboration boosted their experimental sensitivity to (e) neutrinoless double-beta decays through (e&eb) a combination of factors.** |
| **(14) The researchers deployed clean detectors with (e&eb) very low background noise.** |
| **(15) And they used an unprecedented amount of xenon-136, which they purified to (e) remove any radioactive contaminants that could produce (eb) unwanted signals in the detectors.** |
| **(16) Combining data from several years, they improved (eb) the limit on the probability of neutrinoless double-beta decay by (e) sixfold compared to previous searches.** |
| **(17) This probability corresponds to (e&eb) a decay lifetime whose staggering value exceeds (e) 1026 years.** |
| **(18) This measurement allowed (eb) the researchers to set (eb) the most stringent upper limit on (e&eb) the possible mass of Majorana neutrinos (the particle must be lighter than 61–165 meV). This research is published in Physical Review Letters. –Matteo Rini Matteo Rini is the Deputy Editor of Physics.** |
| **How to Reach a Thousand-Second in-Plane Polarization Lifetime with 0.97−GeV/c Deuterons in a Storage Ring G. Guidoboni et al. (JEDI Collaboration) Phys. Rev. Lett. 117, 054801 – Published 28 July 2016 Physics logo See Synopsis: Deuterons Spin together for a Thousand Seconds** |
| **(19) Authors observe a deuteron beam polarization lifetime near 1000 s in (eb) the horizontal plane of (e) a magnetic storage ring (COSY).** |
| **(20) This long spin coherence time is maintained through (e&eb) a combination of beam bunching, electron cooling, (e&eb) sextupole field corrections, and (e&eb) the suppression of collective effects through (e&eb) beam current limits.** |
| **(21) This record lifetime is required for (e) a storage ring search for (e) an intrinsic electric dipole moment on (e&eb) the deuteron at a statistical sensitivity level approaching 10−29 e cm.** |
| **Synopsis: Deuterons Spin together for a Thousand Seconds July 28, 2016 Researchers set a new record for the in-plane spin-alignment lifetime of deuterons circulating in a magnetic storage ring. Synopsis figure E. Stephenson/Indiana University Center for Spacetime Symmetries Sime sentences and figures are deleted due to spatial restrictions.** |
| **(22) The electric dipole moment (EDM) of (e) a fundamental particle or a simple nucleus is believed to be (=) extremely small.** |
| **(23) One might therefore think that it isn't worth measuring this quantity. But it is. A nonzero EDM would signal (eb) a broken symmetry of (e) nature known as CP violation, which is required to (e) explain why the Universe has (e) more matter than antimatter.** |
| **(24) Edward Stephenson from Indiana University, Bloomington, and colleagues has now made a significant stride towards determining (eb) the EDM of the deuteron using (e) a deuteron beam circulating in (e&eb) a magnetic storage ring.** |

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| **(25) They have been able to preserve the alignment of the deuterons' spins for (e) a chunk of time that may be long enough to determine (eb) the deuteron's EDM.** |
| **(26) In experiments of this kind, the EDM can be determined by (e) measuring the rotation of (e) the initially aligned spins of (e) the beam particles as they travel around (e&eb) the ring.** |
| **(27) To obtain a measurement sufficiently accurate to explore (e&eb) CP violation predicted by (e) theories beyond the standard model, the spin alignment has to be (=) preserved for about 1000 s.** |
| **(28) But changes in particle speed—associated with (e&eb) position or direction variations among (e&eb) the particles in the beam—lead to (eb) a rapid loss of the alignment in the plane of the ring.** |
| **(29) Stephenson and co-workers have managed to retain (eb)the alignment of deuterons for the required 1000 s, which is 3 orders of magnitude longer (eb) than the previous record (obtained for an electron beam).** |
| **(30) Their success resulted from (e) capturing the beam into a tight bunch to reduce (e) position and direction spread and from (e) shaping the ring's magnetic field to(e) ensure all of the deuteron orbits around (e&eb) the ring have the same length, despite (e) direction variations. This research is published in Physical Review Letters. –Ana Lopes Ana Lopes is a Senior Editor of Physics.** |
| **Cotunneling Drag Effect in Coulomb-Coupled Quantum Dots A. J. Keller, J. S. Lim, David Sánchez, Rosa López, S. Amasha, J. A. Katine, Hadas Shtrikman, and D. Goldhaber-Gordon Phys. Rev. Lett. 117, 066602 – Published 4 August 2016 Physics logo See Synopsis: Coulomb Drag in a Double Dot** |
| **(31) In Coulomb drag, a current flowing in one conductor can induce (eb) a voltage across an adjacent conductor via (e&eb) the Coulomb interaction.** |
| **(32) The mechanisms yielding drag effects are not always understood, even though drag effects are (=) sufficiently general to be (=) seen in many low-dimensional systems.** |
| **(33) In this Letter, authors observe Coulomb drag in (eb) a Coulomb-coupled double quantum dot and, through (e&eb) both experimental and theoretical arguments, identify cotunneling as (=) essential to obtaining a correct qualitative understanding of the drag behavior.** |
| **Synopsis: Coulomb Drag in a Double Dot August 4, 2016 Electric current passing through a quantum dot can generate current in a nearby dot through a coordinated tunneling mechanism. Synopsis figureA. J. Keller et al., Phys. Rev. Lett. (2016)** |
| **(34) Electrons moving through (e&eb) a conductor can induce (eb) the movement of electrons of a nearby, but isolated, conductor.** |
| **(35) This so-called Coulomb drag is well understood in 2D layered materials, but the theory for nanowires and quantum dots does not (e) fully explain experiments.** |
| **(36) A new study of Coulomb drag in two nearby quantum dots shows (eb) that simultaneous tunneling of electrons—called cotunneling—can help explain the low-temperature behavior of (e) the system.** |

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| **(37) Electric current flows through (e&eb) a quantum dot by the tunneling of (e) electrons.** |
| **(38) When two dots are placed next to each other, the active, or lead, dot can drag current through (e&eb) the passive dot by (e) a process called sequential tunneling.** |
| **(39) In this case, an electron first tunnels onto (e&eb) the active dot, increasing (eb) the probability of tunneling onto (e&eb) the passive dot.** |
| **(40) Thus the passive dot follows (e&eb) the active dot's lead, but always one step behind.** |
| **(41) Sequential tunneling explains (eb) Coulomb drag observations at (eb) high temperature, but no studies have investigated cold temperatures where (e) quantum coherence effects may play a role.** |
| **(42) Andrew Keller from the California Institute of Technology, Pasadena, and his colleagues explored Coulomb drag at (eb) sub-Kelvin temperatures in a pair of quantum dots coupled by (e) capacitance.** |
| **(43) The team applied a current-driving voltage across (e&eb) one dot and detected a drag current in (eb) the other dot, even in regimes where (e) the driving voltage was very small.** |
| **(44) This conflicts with (e&eb) the sequential tunneling model, which predicts (eb) a threshold driving voltage below which (e) the drag current should be zero.** |
| **(45) The team showed that the inclusion of cotunneling into (eb) the drag theory explains the observations.** |
| **(46) Cotunneling could be important to efforts aimed at harvesting a unidirectional current from (e) a fluctuating voltage supply. This research is published in Physical Review Letters. –Michael Schirber Michael Schirber is a Corresponding Editor for Physics based in Lyon, France.** |
| **Maybe it wasn't the Higgs particle after all November 7, 2014 Maybe it wasn't the Higgs particle after all Associate Professor Mads Toudal Frandsen, University of Southern Denmark Credit: University of Southern Denmark** |
| **(47) Last year CERN announced the finding of a new elementary particle, the Higgs particle. But maybe it wasn't the Higgs particle, maybe it just looks like it. And maybe it is not alone. Many calculations indicate that the particle discovered last year in the CERN particle accelerator was (=) indeed the famous Higgs particle.** |
| **(48) Physicists agree that the CERN experiments did find a new particle that had never been seen before, but according to (e&eb) an international research team, there is (=) no conclusive evidence that the particle was indeed the Higgs particle.The research team has scrutinized the existing scientific data from CERN about the newfound particle and published their analysis in the journal Physical Review D. A member of this team is Mads Toudal Frandsen, associate professor at the Center for Cosmology and Particle Physics Phenomenology, Department of Physics, Chemistry and Pharmacy at the University of Southern Denmark. "The CERN data is generally taken as (=) evidence that the particle is the Higgs particle.** |

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| **(49) It is true that the Higgs particle can explain (eb) the data but there can be other explanations, we would also get this data from other particles", Mads Toudal Frandsen explains.** |
| **(50) The researchers' analysis does not (e) debunk the possibility that CERN has discovered the Higgs particle. That is still possible - but it is equally possible that it is a different kind of particle. "The current data is not precise enough to determine exactly what the particle is. It could be a number of other known particles", says Mads Toudal Frandsen. But if it wasn't the Higgs particle, which was found in CERN's particle accelerator, then what was it? "We believe that it may be a so-called techni-higgs particle. This particle is in some ways similar to the Higgs particle - hence half of the name", says Mads Toudal Frandsen.** |
| **(51) Although the techni-higgs particle and Higgs particle can easily be confused in experiments, they are two very different particles belonging to (e) two very different theories of how the universe was created.** |
| **(52) The Higgs particle is the missing piece in (eb) the theory called the Standard Model.** |
| **(53) This theory describes (eb) three of the four forces of nature. But it does not (e) explain what dark matter is - the substance that makes up (eb) most of the universe. A techni-higgs particle, if it exists, is (=) a completely different thing:** |
| **(54) "A techni-higgs particle is not (e) an elementary particle.** |
| **(55) Instead, it consists of so-called techni-quarks, which we believe are elementary. Techni-quarks may bind together in various ways to form (eb) for instance techni-higgs particles, while other combinations may form (eb) dark matter. Authors therefore expect to find several different particles at the LHC, all built by (e) techni-quarks", says Mads Toudal Frandsen.** |
| **(56) If techni-quarks exist, (eb) there must be a force to bind them together so that (e) they can form particles.** |
| **(57) None of the four known forces of nature (gravity, the electromagnetic force, the weak nuclear force and the strong nuclear force) are (=) any good at binding techni-quarks together.** |
| **(58) There must therefore be (=) a yet undiscovered force of nature. This force is called the the Technicolor force.** |
| **(59) What was found last year in CERN's accelerator could thus be (=) either the Higgs particle of the Standard Model or a light techni-higgs particle, composed of (e) two techni-quarks.** |
| **(60) Mads Toudal Frandsen believes that more data from CERN will probably be able to determine (eb) if it was a Higgs or a techni-higgs particle. If CERN gets an even more powerful accelerator, it will in principle be able to observe techni-quarks directly. More information: Technicolor Higgs boson in the light of LHC data. Phys. Rev. D 90, 035012th Alexander Belyaev, Matthew S. Brown, Roshan Foadi, and Mads T. Frandsen. journals.aps.org/prd/abstract/10.1103/PhysRevD.90.035012. On Arxiv: arxiv.org/abs/1309.2097 Journal reference: Physical Review D search and more info website: Provided by: University of Southern Denmark search and more info website** |

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| **The inflationary brane-antibrane universeClifford P. Burgess1,2, Mahbub Majumdar3, Detlef Nolte1, Fernando Quevedo3, Govindan Rajesh1 and Ren-Jie Zhang1 Published 13 August 2001 • Journal of High Energy Physics, Volume 2001, JHEP07(2001)** |
| **(61) Authors show how the motion through (e&eb) the extra dimensions of a gas of branes and antibranes can, under certain circumstances, produce (eb) an era of inflation as seen by (e) observers trapped on (e&eb) a 3-brane, with (e&eb) the inflaton being the inter-brane separation.** |
| **(62) Although most of discussion refers to (e) arbitrary ps, when we need to be specific they assume that they are (=) D-branes of type-II or type-I string theory.** |
| **(63) For realistic brane couplings, such as those arising in (eb) string theory, the inter-brane potentials are (=) too steep to inflate (eb) the universe for acceptably long times.** |
| **(64) However, for special regions of the parameter space of (e) brane-antibrane positions the brane motion is (=) slow enough for there to be (=) sufficient inflation.** |
| **(65) Inflation would be more generic in models where (e) the inter-brane interactions are much weaker.** |
| **(66) The spectrum of primordial density fluctuations predicted has (e) index n slightly less than 1, and acceptable amplitude, provided (e) that the extra dimensions have (e) linear size 1/r ~ 1012 GeV.** |
| **(67) Reheating occurs as in hybrid inflation, with (e&eb) the tachyonic instability of (e) the brane-antibrane system taking over for (e) small separations.** |
| **(68) The tachyon field can induce (eb) a cascade mechanism within which higher-dimension branes annihilate (e) into lower-dimension ones.** |
| **(69) They argue that such a cascade naturally stops with (e&eb) the production of 3-branes in (eb) 10-dimensional string theory.** |
| **Effect of the neutron lifetime on processes in the early universe A. V. Chechkin, A. V. Ivanchik, A. P. Serebrov, S. V. Bobashev First Online: 27 July 2016 DOI: 10.1134/S1063784216070069 Chechkin, A.V., Ivanchik, A.V., Serebrov, A.P. et al. Tech. Phys. (2016) 61: 1101. Doi: 10.1134/S1063784216070069** |
| **(70) The influence of the neutron lifetime on (e&eb) the abundance of light elements produced during (e&eb) the primordial nucleosynthesis early in (eb) the birth of the Universe is considered.** |
| **(71) Among light elements, namely, D, 3He, 4He, and 7Li, 4He proves to be (=) most sensitive to neutron lifetime τn.** |
| **(72) Astronomic data on the light element abundance also provide (eb) the best accuracy for 4He.** |
| **(73) The solution of a number of problems discussed in this paper requires (e) improving the accuracy of observations for the 4He abundance and refining the value of τn.** |

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| **Application of a Theory and Simulation based Convective Boundary Mixing model for AGB Star Evolution and Nucleosynthesis U. Battino, M. Pignatari, C. Ritter, F. Herwig, P. Denisenkov, J.W. Den Hartogh, R. Trappitsch, R. Hirschi, B. Freytag, F. Thielemann, B. Paxton (Submitted on 19 May 2016)** |
| **(74) The s-process nucleosynthesis in Asymptotic Giant Branch (AGB) stars depends on (e&eb) the modeling of (e) convective boundaries.** |
| **(75) Authors' present models and s-process simulations that adopt (e&eb) a treatment of convective boundaries based on (e) the results of hydrodynamic simulations and on (e&eb) the theory of mixing due to (e) gravity waves in (eb) the vicinity of convective boundaries.** |
| **(76) Hydrodynamics simulations suggest (eb) the presence of convective boundary mixing (CBM) at (eb) the bottom of the thermal pulse-driven convective zone.** |
| **(77) Similarly, convection-induced mixing processes are proposed for (e) the mixing below the convective envelope during (e&eb) third dredge-up where (e) the 13C pocket for (e) the s process in AGB stars forms.** |
| **(78) In this work apply apply a CBM model motivated by (e) simulations and theory to (e&eb) models with initial mass M = 2 and M = 3M?, and with (e&eb) initial metal content Z = 0:01 and Z = 0:02.** |
| **(79) As reported previously, the He-intershell abundance of 12C and 16O are increased by (e) CBM at the bottom of (e) pulse-driven convection zone.** |
| **(80) This mixing is affecting (e&eb) the 22Ne (alpha, n) 25Mg activation and (e&eb) the s-process effciency in the 13C-pocket.** |
| **(81) In our model CBM at the bottom of the convective envelope during (e&eb) the third dredgeup represents (eb) gravity wave mixing.** |
| **(82) Authors take further into account that hydrodynamic simulations indicate (eb) a declining mixing efficiency already about (e&eb) a pressure scale height from (e) the convective boundaries, compared to (e&eb) mixing-length theory.** |
| **(83) The final s-process abundances are characterized by (e) 0:36 < [s=Fe] < 0:78 and the heavy-to-light s-process ratio is (=) 0:23 < [hs=ls] < 0:45.** |
| **(84) Finally, authors compare results with (e&eb) stellar observations, pre-solar grain measurements and previous work. Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI:10.3847/0004-637X/827/1/30 Cite as: arXiv:1605.06159 [astro-ph.SR] (or arXiv:1605.06159v1 [astro-ph.SR] for this version)** |
| **MESA and NuGrid simulations of classical novae: CO and ONe nova nucleosynthesis Pavel A. Denissenkov (UVic), James W. Truran (Chicago), Marco Pignatari (Basel), Reto Trappitsch (Chicago), Christian Ritter (Frankfurt), Falk Herwig (UVic), Umberto Battino (Basel), Kiana Setoodehnia (McMaster), B. Paxton (UC Santa Barbara)** |

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| **(85) Classical novae are the result of (e) thermonuclear flashes of hydrogen accreted by (e) CO or ONe white dwarfs, leading (eb) eventually to the dynamic ejection of (e) the surface layers.** |
| **(86) These are observationally known to be (=) enriched in heavy elements, such as (=) C, O and Ne that must originate in (eb) layers below the H-flash convection zone.** |
| **(87) Building on their previous work, they now present stellar evolution simulations of (e) ONe novae and provide (eb) a comprehensive comparison of models with (e&eb) published ones.** |
| **(88) Some of authors models include exponential convective boundary mixing to account for (e) the observed enrichment of (e) the nova ejecta even when accreted material has (e) a solar abundance distribution.** |
| **(89) Models produce (eb) maximum temperature evolution profiles and nucleosynthesis yields (eb) in good agreement with (e&eb) models that generate (eb) enriched ejecta by assuming (eb) that the accreted material was pre-mixed.** |
| **(90) Authors confirm for (e) ONe novae the result we reported previously, i.e.\ we found that 3He could be produced {\it in situ} in (eb) solar-composition envelopes accreted with (e&eb) slow rates ($M˙<10−10M⊙$/yr) by (e) cold (TWD<107 K) CO WDs, and that convection was triggered by (eb) 3He burning before (e) the nova outburst in that case.** |
| **(91) In addition, authors now find that the interplay between the 3He production and (e&eb) destruction in the solar-composition envelope accreted with (e&eb) an intermediate rate, e.g.\ $M˙=10−10M⊙$/yr, by the $1.15M⊙$ ONe WD with (e&eb) a relatively high initial central temperature, e.g.\ TWD=15×106 K, leads to (eb) the formation of a thick radiative buffer zone that separates (e&eb) the bottom of the convective envelope from (e) the WD surface. (Abridged) Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI: 10.1093/mnras/stu1000 Cite as: arXiv:1303.6265 [astro-ph.SR] (or arXiv:1303.6265v2 [astro-ph.SR] for this version)** |
| **Hybrid C-O-Ne White Dwarfs as Progenitors of Diverse SNe Ia Pavel Denissenkov (UVic), James Truran (U. Chicago), Falk Herwig (UVic), Sam Jones (UVic), Bill Paxton (UC Santa Barbara), Ken'ichi Nomoto (U. Tokyo), Toshio Suzuki (Nihon U.), Hiroshi Toki (Osaka U.)(Submitted on 6 Nov 2014)** |
| **(92) When carbon is ignited off-center in a CO core of a super-AGB star (eb) its burning in a convective shell tends to propagate to (e&eb) the center.** |
| **(93) Whether the C flame will actually be able to reach (eb) the center depends on (e&eb) the efficiency of (e) extra mixing beneath (e&eb) the C convective shell.** |
| **(94) Whereas thermohaline mixing is too inefficient to interfere with the C-flame propagation, (eb) convective boundary mixing can prevent (e) the C burning from reaching (eb) the center.** |
| **(95) As a result, a C-O-Ne white dwarf (WD) is (=) formed, after the star has lost (e) its envelope.** |
| **(96) Such a "hybrid" WD has (e) a small CO core surrounded by (e) a thick ONe zone.** |

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| **(97) In the proposed 1D stellar evolution computations the hybrid WD is allowed to (eb) accrete C-rich material, as if it were in a close binary system and (e&eb) accreted H-rich material from (e) its companion with (e&eb) a sufficiently high rate at which the accreted H would be processed into (e&eb) He under (e&eb) stationary conditions, assuming (eb) that He could then be transformed into (e&eb) C.** |
| **(98) When the mass of the accreting WD approaches (eb) the Chandrasekhar limit, they find a series of convective Urca shell flashes associated with (e&eb) high abundances of 23Na and 25Mg.** |
| **(99) They are followed by off-center C ignition leading to (eb) convection that occupies (e) almost the entire star.** |
| **(100) To model the Urca processes, authors use the most recent well-resolved data for (e) their reaction and (e&eb) neutrino-energy loss rates.** |
| **(101) Because of the emphasized uncertainty of the convective Urca process in hybrid WD models of (e) SN Ia progenitors, authors consider a number of their potentially possible alternative instances for (e) different mixing assumptions, all of which reach (eb) a phase of explosive C ignition, either off or in (eb) the center.** |
| **(102) Our hybrid SN Ia progenitor models have (e) much lower C to O abundance ratios at (eb) the moment of the explosive C ignition than (e) their pure CO counterparts, which may explain (eb) the observed diversity of the SNe Ia. subjects: Solar and Stellar Astrophysics (astro-ph.SR)cite as: arXiv:1411.1471 [astro-ph.SR](or arXiv:1411.1471v1 [astro-ph.SR] for this version)** |
| **Rejecting proposed dense-matter equations of state with quiescent low-mass X-ray binaries Sebastien Guillot, Robert E. Rutledge (Submitted on 15 Sep 2014 (v1), last revised 20 Oct 2014 (this version, v2))** |
| **(103) Neutrons stars are unique laboratories to (e) discriminate between the various proposed equations of (e) state of matter at and above nuclear density.** |
| **(104) One sub-class of neutron stars - those inside (eb) quiescent low-mass X-ray binaries (qLMXBs) - produce (eb) a thermal surface emission from (e) which the neutron star radius (R_NS) can be measured, using (e) the widely accepted observational scenario for qLMXBs, assuming (eb) unmagnetized H atmospheres.** |
| **(105) In a combined spectral analysis, this work first reproduces (eb) a previously published measurement of the \rns, assumed to be (=) the same for all neutron stars, using (e) a slightly expanded data set.** |
| **(106) The radius measured is R_NS = 9.4 +/-1.2 km. On the basis of spectral analysis alone, this measured value is not affected by imposing (e&eb) an assumption of causality in the core.** |
| **(107) However, the assumptions underlying this R_NS measurement would be falsified by (e) the observation of any neutron star with (e&eb) a mass >2.6 Msun, since radii <11 km would be rejected if (e) causality is assumed, which would exclude (eb) most of the R_NS parameter space obtained in this analysis.** |

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| **(108) Finally, this work directly tests (e&eb) a selection of dense matter equations of states: WFF1, AP4, MPA1, PAL1, MS0, and three versions of equations of state produced through (e&eb) chiral effective theory.** |
| **(109) Two of those, MSO and PAL1, are rejected at (eb) the 99% confidence level, accounting for (e) all quantifiable uncertainties, while (e) the other cannot be excluded at >99% certainty. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE) DOI: 10.1088/2041-8205/796/1/L3 Cite as: arXiv:1409.4306 [astro-ph.HE] (or arXiv:1409.4306v2 [astro-ph.HE] for this version** |

## SECTION SEVEN

## Formula Should Be E= (E0)/(1-V^2/C^2)^(1/2) Derived By Laue In 1911 And Klein In 1918

**INTRODUCTION—VARIABLES USED**

**Is E=mc^2 an exclusively relativistic result? Qing-Ping Ma**

(1) The derivation by Einstein in 1905 is logically flawed as (e) a relativistic proof and the truly relativistic formula should be E= (E0)/(1-v^2/c^2)^(1/2) derived by Laue in 1911 and Klein in 1918.
(2) If the energy measured in one reference frame is E0, it is E= (E0)/(1-v^2/c^2)^(1/2) in a reference frame moving at velocity v relative to (e&eb) the first frame. Subjects: General Physics (physics.gen-ph); History and Philosophy of Physics (physics.hist-ph) Cite as: arXiv:1509.07517 [physics.gen-ph] (or arXiv:1509.07517v1 [physics.gen-ph] for this version)

**GW150914: Implications for the Stochastic Gravitational-Wave Background from Binary Black Holes B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 131102 – Published 31 March 2016**

(1) The LIGO detection of the gravitational wave transient GW150914, from (e) the inspiral and merger of (e) two black holes with masses ≳30M⊙, suggests (eb) a population of binary black holes with (e&eb) relatively high mass.
(2) This observation implies (eb) that the stochastic gravitational-wave background from binary black holes, created from (e) the incoherent superposition of (e) all the merging binaries in the Universe, could be (=) higher than previously expected.
(3) Using the properties of GW150914, authors estimate the energy density of (e) such a background from (e) binary black holes.
(4) In the most sensitive part of the Advanced LIGO and Advanced Virgo band for (e) stochastic backgrounds (near 25 Hz), authors predict ΩGW(f=25 Hz)=1.1+2.7−0.9×10−9 with 90% confidence.
(5) This prediction is robustly demonstrated for (e) a variety of formation scenarios with (e&eb) different parameters.
(6) The differences between models are (=) small compared to the statistical uncertainty arising from

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| (e) the currently poorly constrained (e) local coalescence rate.<br>(7) Authors conclude that this background is potentially measurable by (e) the Advanced LIGO and Advanced Virgo detectors operating at their projected final sensitivity. DOI:http://dx.doi.org/10.1103/PhysRevLett.116.131102 © 2016 American Physical Society | |
| **NOTATION** | |
| **Module One**<br>The derivation by Einstein in 1905 is logically flawed as (e) a relativistic proof and the truly relativistic formula should be E= (E0)/(1-v^2/c^2)^(1/2) derived by Laue in 1911 and Klein in 1918 | |
| $G_{13}$ : Category one of relativistic proof and the truly relativistic formula should be E= (E0)/(1-v^2/c^2)^(1/2) derived by Laue in 1911 and Klein in 1918<br>$G_{14}$ : Category two of SAS<br>$G_{15}$ : Category three of SAS | |
| $T_{13}$ : Category one of derivation by Einstein in 1905 is logically flawed<br>$T_{14}$ : Category two of SAS<br>$T_{15}$ : Category three of SAS | |
| **Module Two**<br>E= (E0)/(1-v^2/c^2)^(1/2) | |
| $G_{16}$ : Category one of LHS of the above equation<br>$G_{17}$ : Category two of SAS<br>$G_{18}$ : Category three of SAS | |
| $T_{16}$ : Category one of RHS of the above equation<br>$T_{17}$ : Category two of SAS<br>$T_{18}$ : Category three of SAS | |
| **Module three**<br>If the energy measured in one reference frame is E0, it is E= (E0)/(1-v^2/c^2)^(1/2) in a reference frame moving at velocity v relative to (e&eb) the first frame.<br>Subjects: General Physics (physics.gen-ph); History and Philosophy of Physics (physics.hist-ph)<br>Cite as: arXiv:1509.07517 [physics.gen-ph] (or arXiv:1509.07517v1 [physics.gen-ph] for this version | |
| $G_{20}$ : Category one of energy measured in one reference frame is E0, | |

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| $G_{21}$ : Category two of SAS<br><br>$G_{22}$ : Category three of SAS | |
| $T_{20}$ : Category one of E= (E0)/(1-v^2/c^2)^(1/2) in a reference frame moving at velocity v relative to (e&eb) the first frame.<br><br>$T_{21}$ : Category two of SAS<br><br>$T_{22}$ : Category three of SAS | |
| **Module four**<br><br>If the energy measured in one reference frame is E0, it is E= (E0)/ (1-v^2/c^2) ^ (1/2) in a reference frame moving at velocity v relative to (e&eb) the first frame.<br><br>Subjects: General Physics (physics.gen-ph); History and Philosophy of Physics (physics.hist-ph)<br>Cite as: arXiv:1509.07517 [physics.gen-ph] (or arXiv:1509.07517v1 [physics.gen-ph] for this version | |
| $G_{24}$ : Category one of **energy measured in one reference frame is E0, it is E= (E0)/ (1-v^2/c^2) ^ (1/2) in a reference frame moving at velocity v relative;** first frame.<br><br>$G_{25}$ : Category two of SAS<br><br>$G_{26}$ : Category three of SAS | |
| $T_{24}$ : Category one of first frame.; **energy measured in one reference frame is E0, it is E= (E0)/ (1-v^2/c^2) ^ (1/2) in a reference frame moving at velocity v relative**<br><br>$T_{25}$ : Category two of SAS<br><br>$T_{26}$ : Category three of SAS | |
| **Module five**<br><br>The LIGO detection of the gravitational wave transient GW150914, from (e) the inspiral and merger of (e) two black holes with masses ≳30M⊙, suggests (eb) a population of binary black holes with (e&eb) relatively high mass | |
| $G_{28}$ : Category one of inspiral and merger of (e) two black holes with masses ≳30M⊙, suggests (eb) a population of binary black holes with (e&eb) relatively high mass<br><br>$G_{29}$ : Category two of SAS<br><br>$G_{30}$ : Category three of SAS | |
| $T_{28}$ : Category one of LIGO detection of the gravitational wave transient GW150914<br><br>$T_{29}$ : Category two of SAS | |

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| $T_{30}$ : Category three of SAS | |
| **Module six**<br><br>The LIGO detection of the gravitational wave transient GW150914, from (e) the inspiral and merger of two black holes with masses ≳30M⊙, suggests (eb) a population of binary black holes with (e&eb) relatively high mass | |
| $G_{32}$ : Category one of LIGO detection of the gravitational wave transient GW150914, from (e) the inspiral and merger of two black holes with masses ≳30M⊙,<br><br>$G_{33}$ : Category two of SAS<br><br>$G_{34}$ : Category three of SAS | |
| $T_{32}$ : Category one of population of binary black holes with (e&eb) relatively high mass<br><br>$T_{33}$ : Category two of SAS<br><br>$T_{34}$ : Category three of SAS | |
| **Module seven**<br><br>The LIGO detection of the gravitational wave transient GW150914, from the inspiral and merger of two black holes with masses ≳30M⊙, suggests (eb) a population of binary black holes with (e&eb) relatively high mass | |
| $G_{36}$ : Category one of LIGO detection of the gravitational wave transient GW150914, from the inspiral and merger of two black holes with masses ≳30M⊙,<br><br>$G_{37}$ : Category two of SAS<br><br>$G_{38}$ : Category three of SAS | |
| $T_{36}$ : Category one of population of binary black holes with (e&eb) relatively high mass<br><br>$T_{37}$ : Category two of SAS<br><br>$T_{38}$ : Category three of SAS | |
| **Module eight**<br><br>The LIGO detection of the gravitational wave transient GW150914, from the inspiral and merger of two black holes with masses ≳30M⊙, suggests a population of binary black holes with (e&eb) relatively high mass | |
| $G_{40}$ : Category one of **LIGO detection of the gravitational wave transient GW150914, from the inspiral and merger of two black holes with masses ≳30M⊙, suggests a population of binary black holes;** relatively high mass | |

$G_{41}$ : Category two of SAS

$G_{42}$ : Category three of SAS

$T_{40}$ : Category one of relatively high mass **;LIGO detection of the gravitational wave transient GW150914, from the inspiral and merger of two black holes with masses ≳30M⊙, suggests a population of binary black holes**

$T_{41}$ : Category two of SAS

$T_{42}$ : Category three of SAS

**Module Nine**

LIGO detection of the gravitational wave transient GW150914, from the inspiral and merger of two black holes with masses ≳30M⊙, implies (eb) that the stochastic gravitational-wave background from binary black holes, created from (e) the incoherent superposition of (e) all the merging binaries in the Universe, could be (=) higher than previously expected

$G_{44}$ : Category one of LIGO detection of the gravitational wave transient GW150914, from the inspiral and merger of two black holes with masses ≳30M⊙,

$G_{45}$ : Category two of SAS

$G_{46}$ : Category three of SAS

$T_{44}$ : Category one ofstochastic gravitational-wave background from binary black holes, created from (e) the incoherent superposition of (e) all the merging binaries in the Universe, could be (=) higher than previously expected

$T_{45}$ : Category two of SAS

$T_{46}$ : Category three of SAS

**The Coefficients:**

$(a_{13})^{(1)}, (a_{14})^{(1)}, (a_{15})^{(1)}, (b_{13})^{(1)}, (b_{14})^{(1)}, (b_{15})^{(1)}$ $(a_{16})^{(2)}, (a_{17})^{(2)}, (a_{18})^{(2)}$ $(b_{16})^{(2)}, (b_{17})^{(2)}, (b_{18})^{(2)}$:
$(a_{20})^{(3)}, (a_{21})^{(3)}, (a_{22})^{(3)}$ , $(b_{20})^{(3)}, (b_{21})^{(3)}, (b_{22})^{(3)}$
$(a_{24})^{(4)}, (a_{25})^{(4)}, (a_{26})^{(4)}, (b_{24})^{(4)}, (b_{25})^{(4)}, (b_{26})^{(4)}, (b_{28})^{(5)}, (b_{29})^{(5)}, (b_{30})^{(5)}, (a_{28})^{(5)}, (a_{29})^{(5)}, (a_{30})^{(5)}$,
$(a_{32})^{(6)}, (a_{33})^{(6)}, (a_{34})^{(6)}, (b_{32})^{(6)}, (b_{33})^{(6)}, (b_{34})^{(6)}$
$(a_{36})^{(7)}, (a_{37})^{(7)}, (a_{38})^{(7)}, (b_{36})^{(7)}, (b_{37})^{(7)}, (b_{38})^{(7)}$
$(a_{40})^{(8)}, (a_{41})^{(8)}, (a_{42})^{(8)}, (b_{40})^{(8)}, (b_{41})^{(8)}, (b_{42})^{(8)}$
$(a_{44})^{(9)}, (a_{45})^{(9)}, (a_{46})^{(9)}, (b_{44})^{(9)}, (b_{45})^{(9)}, (b_{46})^{(9)}$

are Accentuation coefficients
$(a'_{13})^{(1)}, (a'_{14})^{(1)}, (a'_{15})^{(1)}, (b'_{13})^{(1)}, (b'_{14})^{(1)}, (b'_{15})^{(1)}, (a'_{16})^{(2)}, (a'_{17})^{(2)}, (a'_{18})^{(2)}$, $(b'_{16})^{(2)}, (b'_{17})^{(2)}, (b'_{18})^{(2)}$
$, (a'_{20})^{(3)}, (a'_{21})^{(3)}, (a'_{22})^{(3)}, (b'_{20})^{(3)}, (b'_{21})^{(3)}, (b'_{22})^{(3)}$
$(a'_{24})^{(4)}, (a'_{25})^{(4)}, (a'_{26})^{(4)}, (b'_{24})^{(4)}, (b'_{25})^{(4)}, (b'_{26})^{(4)}, (b'_{28})^{(5)}, (b'_{29})^{(5)}, (b'_{30})^{(5)}$ $(a'_{28})^{(5)}, (a'_{29})^{(5)}, (a'_{30})^{(5)}$ ,
$(a'_{32})^{(6)}, (a'_{33})^{(6)}, (a'_{34})^{(6)}, (b'_{32})^{(6)}, (b'_{33})^{(6)}, (b'_{34})^{(6)}$

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| $(a_{36}^{'})^{(7)}, (a_{37}^{'})^{(7)}, (a_{38}^{'})^{(7)}, (b_{36}^{'})^{(7)}, (b_{37}^{'})^{(7)}, (b_{38}^{'})^{(7)},$ $(a_{40}^{'})^{(8)}, (a_{41}^{'})^{(8)}, (a_{42}^{'})^{(8)}, (b_{40}^{'})^{(8)}, (b_{41}^{'})^{(8)}, (b_{42}^{'})^{(8)},$ $(a_{44}^{'})^{(9)}, (a_{45}^{'})^{(9)}, (a_{46}^{'})^{(9)}, (b_{44}^{'})^{(9)}, (b_{45}^{'})^{(9)}, (b_{46}^{'})^{(9)},$ are Dissipation coefficients | |
| **Module Numbered One**<br><br>**The differential system of this model is now (Module Numbered one)** | |
| $\frac{dG_{13}}{dt} = (a_{13})^{(1)}G_{14} - \left[(a_{13}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14}, t)\right]G_{13}$ | 1 |
| $\frac{dG_{14}}{dt} = (a_{14})^{(1)}G_{13} - \left[(a_{14}^{'})^{(1)} + (a_{14}^{''})^{(1)}(T_{14}, t)\right]G_{14}$ | 2 |
| $\frac{dG_{15}}{dt} = (a_{15})^{(1)}G_{14} - \left[(a_{15}^{'})^{(1)} + (a_{15}^{''})^{(1)}(T_{14}, t)\right]G_{15}$ | 3 |
| $\frac{dT_{13}}{dt} = (b_{13})^{(1)}T_{14} - \left[(b_{13}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G, t)\right]T_{13}$ | 4 |
| $\frac{dT_{14}}{dt} = (b_{14})^{(1)}T_{13} - \left[(b_{14}^{'})^{(1)} - (b_{14}^{''})^{(1)}(G, t)\right]T_{14}$ | 5 |
| $\frac{dT_{15}}{dt} = (b_{15})^{(1)}T_{14} - \left[(b_{15}^{'})^{(1)} - (b_{15}^{''})^{(1)}(G, t)\right]T_{15}$ | 6 |
| $+(a_{13}^{''})^{(1)}(T_{14}, t) =$ First augmentation factor | |
| $-(b_{13}^{''})^{(1)}(G, t) =$ First detritions factor | |
| **Module Numbered Two**<br><br>**The differential system of this model is now ( Module numbered two)** | |
| $\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \left[(a_{16}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17}, t)\right]G_{16}$ | 7 |
| $\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \left[(a_{17}^{'})^{(2)} + (a_{17}^{''})^{(2)}(T_{17}, t)\right]G_{17}$ | 8 |
| $\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \left[(a_{18}^{'})^{(2)} + (a_{18}^{''})^{(2)}(T_{17}, t)\right]G_{18}$ | 9 |
| $\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \left[(b_{16}^{'})^{(2)} - (b_{16}^{''})^{(2)}\left((G_{19}), t\right)\right]T_{16}$ | 10 |
| $\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \left[(b_{17}^{'})^{(2)} - (b_{17}^{''})^{(2)}\left((G_{19}), t\right)\right]T_{17}$ | 11 |
| $\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \left[(b_{18}^{'})^{(2)} - (b_{18}^{''})^{(2)}\left((G_{19}), t\right)\right]T_{18}$ | 12 |
| $+(a_{16}^{''})^{(2)}(T_{17}, t) =$ First augmentation factor | |
| $-(b_{16}^{''})^{(2)}\left((G_{19}), t\right) =$ First detritions factor | |
| **Module Numbered Three**<br><br>**The differential system of this model is now (Module numbered three)** | |
| $\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \left[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21}, t)\right]G_{20}$ | 13 |
| $\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \left[(a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)}(T_{21}, t)\right]G_{21}$ | 14 |
| $\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \left[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21}, t)\right]G_{22}$ | 15 |
| $\frac{dT_{20}}{dt} = (b_{20})^{(3)}T_{21} - \left[(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23}, t)\right]T_{20}$ | 16 |
| $\frac{dT_{21}}{dt} = (b_{21})^{(3)}T_{20} - \left[(b_{21}^{'})^{(3)} - (b_{21}^{''})^{(3)}(G_{23}, t)\right]T_{21}$ | 17 |
| $\frac{dT_{22}}{dt} = (b_{22})^{(3)}T_{21} - \left[(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23}, t)\right]T_{22}$ | 18 |
| $+(a_{20}^{''})^{(3)}(T_{21}, t) =$ First augmentation factor | |
| $-(b_{20}^{''})^{(3)}(G_{23}, t) =$ First detritions factor | |

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| **Module Numbered Four**<br><br>**The differential system of this model is now (Module numbered Four)** | |
| $\frac{dG_{24}}{dt} = (a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25}, t)\right]G_{24}$ | 19 |
| $\frac{dG_{25}}{dt} = (a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25}, t)\right]G_{25}$ | 20 |
| $\frac{dG_{26}}{dt} = (a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25}, t)\right]G_{26}$ | 21 |
| $\frac{dT_{24}}{dt} = (b_{24})^{(4)}T_{25} - \left[(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}((G_{27}), t)\right]T_{24}$ | 22 |
| $\frac{dT_{25}}{dt} = (b_{25})^{(4)}T_{24} - \left[(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}((G_{27}), t)\right]T_{25}$ | 23 |
| $\frac{dT_{26}}{dt} = (b_{26})^{(4)}T_{25} - \left[(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}((G_{27}), t)\right]T_{26}$ | 24 |
| $+(a_{24}^{''})^{(4)}(T_{25}, t) =$ **First augmentation factor** | |
| $-(b_{24}^{''})^{(4)}((G_{27}), t) =$ **First detritions factor** | |
| **Module Numbered Five**:<br><br>**The differential system of this model is now (Module number five)** | |
| $\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}, t)\right]G_{28}$ | 25 |
| $\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29}, t)\right]G_{29}$ | 26 |
| $\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \left[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29}, t)\right]G_{30}$ | 27 |
| $\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \left[(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}((G_{31}), t)\right]T_{28}$ | 28 |
| $\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \left[(b_{29}^{'})^{(5)} - (b_{29}^{''})^{(5)}((G_{31}), t)\right]T_{29}$ | 29 |
| $\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \left[(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}((G_{31}), t)\right]T_{30}$ | 30 |
| $+(a_{28}^{''})^{(5)}(T_{29}, t) =$ **First augmentation factor** | |
| $-(b_{28}^{''})^{(5)}((G_{31}), t) =$ **First detritions factor** | |
| **Module Numbered Six**<br><br>**The differential system of this model is now (Module numbered Six)** | |
| $\frac{dG_{32}}{dt} = (a_{32})^{(6)}G_{33} - \left[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33}, t)\right]G_{32}$ | 31 |
| $\frac{dG_{33}}{dt} = (a_{33})^{(6)}G_{32} - \left[(a_{33}^{'})^{(6)} + (a_{33}^{''})^{(6)}(T_{33}, t)\right]G_{33}$ | 32 |
| $\frac{dG_{34}}{dt} = (a_{34})^{(6)}G_{33} - \left[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33}, t)\right]G_{34}$ | 33 |
| $\frac{dT_{32}}{dt} = (b_{32})^{(6)}T_{33} - \left[(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}((G_{35}), t)\right]T_{32}$ | 34 |
| $\frac{dT_{33}}{dt} = (b_{33})^{(6)}T_{32} - \left[(b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}((G_{35}), t)\right]T_{33}$ | 35 |
| $\frac{dT_{34}}{dt} = (b_{34})^{(6)}T_{33} - \left[(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}((G_{35}), t)\right]T_{34}$ | 36 |
| $+(a_{32}^{''})^{(6)}(T_{33}, t) =$ **First augmentation factor** | |
| **Module Numbered Seven:**<br><br>**The differential system of this model is now (Seventh Module)** | |
| $\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \left[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37}, t)\right]G_{36}$ | 37 |
| $\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \left[(a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37}, t)\right]G_{37}$ | *38* |

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| $\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \left[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37},t)\right]G_{38}$ | *39* |
| $\frac{dT_{36}}{dt} = (b_{36})^{(7)}T_{37} - \left[(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}((G_{39}),t)\right]T_{36}$ | *40* |
| $\frac{dT_{37}}{dt} = (b_{37})^{(7)}T_{36} - \left[(b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}((G_{39}),t)\right]T_{37}$ | *41* |
| $\frac{dT_{38}}{dt} = (b_{38})^{(7)}T_{37} - \left[(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}((G_{39}),t)\right]T_{38}$ | *42* |
| $+(a_{36}^{''})^{(7)}(T_{37},t) =$ **First augmentation factor** | |
| **Module Numbered Eight**<br><br>The differential system of this model is now | |
| $\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \left[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41},t)\right]G_{40}$ | *43* |
| $\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \left[(a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41},t)\right]G_{41}$ | *44* |
| $\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \left[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41},t)\right]G_{42}$ | *45* |
| $\frac{dT_{40}}{dt} = (b_{40})^{(8)}T_{41} - \left[(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}((G_{43}),t)\right]T_{40}$ | *46* |
| $\frac{dT_{41}}{dt} = (b_{41})^{(8)}T_{40} - \left[(b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}((G_{43}),t)\right]T_{41}$ | *47* |
| $\frac{dT_{42}}{dt} = (b_{42})^{(8)}T_{41} - \left[(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}((G_{43}),t)\right]T_{42}$ | *48* |
| **Module Numbered Nine**<br><br>The differential system of this model is now | |
| $\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \left[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45},t)\right]G_{44}$ | 49 |
| $\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \left[(a_{45}^{'})^{(9)} + (a_{45}^{''})^{(9)}(T_{45},t)\right]G_{45}$ | 50 |
| $\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \left[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45},t)\right]G_{46}$ | 51 |
| $\frac{dT_{44}}{dt} = (b_{44})^{(9)}T_{45} - \left[(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}((G_{47}),t)\right]T_{44}$ | 52 |
| $\frac{dT_{45}}{dt} = (b_{45})^{(9)}T_{44} - \left[(b_{45}^{'})^{(9)} - (b_{45}^{''})^{(9)}((G_{47}),t)\right]T_{45}$ | 53 |
| $\frac{dT_{46}}{dt} = (b_{46})^{(9)}T_{45} - \left[(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}((G_{47}),t)\right]T_{46}$ | 54 |
| $+(a_{44}^{''})^{(9)}(T_{45},t) =$ **First augmentation factor** | |
| $-(b_{44}^{''})^{(9)}((G_{47}),t) =$ **First detrition factor** | |
| $\frac{dG_{13}}{dt} = (a_{13})^{(1)}G_{14} - \left[\begin{matrix}(a_{13}^{'})^{(1)}\boxed{+(a_{13}^{''})^{(1)}(T_{14},t)}\boxed{+(a_{16}^{''})^{(2,2,)}(T_{17},t)}\boxed{+(a_{20}^{''})^{(3,3,)}(T_{21},t)}\\ \boxed{+(a_{24}^{''})^{(4,4,4,4,)}(T_{25},t)}\boxed{+(a_{28}^{''})^{(5,5,5,5,)}(T_{29},t)}\boxed{+(a_{32}^{''})^{(6,6,6,6,)}(T_{33},t)}\\ \boxed{+(a_{36}^{''})^{(7,7)}(T_{37},t)}\boxed{+(a_{40}^{''})^{(8,8)}(T_{41},t)}\boxed{+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)}\end{matrix}\right]G_{13}$ | *55* |
| $\frac{dG_{14}}{dt} = (a_{14})^{(1)}G_{13} - \left[\begin{matrix}(a_{14}^{'})^{(1)}\boxed{+(a_{14}^{''})^{(1)}(T_{14},t)}\boxed{+(a_{17}^{''})^{(2,2,)}(T_{17},t)}\boxed{+(a_{21}^{''})^{(3,3,)}(T_{21},t)}\\ \boxed{+(a_{25}^{''})^{(4,4,4,4,)}(T_{25},t)}\boxed{+(a_{29}^{''})^{(5,5,5,5,)}(T_{29},t)}\boxed{+(a_{33}^{''})^{(6,6,6,6,)}(T_{33},t)}\\ \boxed{+(a_{37}^{''})^{(7,7)}(T_{37},t)}\boxed{+(a_{41}^{''})^{(8,8)}(T_{41},t)}\boxed{+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)}\end{matrix}\right]G_{14}$ | *56* |
| $\frac{dG_{15}}{dt} = (a_{15})^{(1)}G_{14} - \left[\begin{matrix}(a_{15}^{'})^{(1)}\boxed{+(a_{15}^{''})^{(1)}(T_{14},t)}\boxed{+(a_{18}^{''})^{(2,2,)}(T_{17},t)}\boxed{+(a_{22}^{''})^{(3,3,)}(T_{21},t)}\\ \boxed{+(a_{26}^{''})^{(4,4,4,4,)}(T_{25},t)}\boxed{+(a_{30}^{''})^{(5,5,5,5,)}(T_{29},t)}\boxed{+(a_{34}^{''})^{(6,6,6,6,)}(T_{33},t)}\\ \boxed{+(a_{38}^{''})^{(7,7)}(T_{37},t)}\boxed{+(a_{42}^{''})^{(8,8)}(T_{41},t)}\boxed{+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)}\end{matrix}\right]G_{15}$ | *57* |
| Where $\boxed{(a_{13}^{''})^{(1)}(T_{14},t)}$, $\boxed{(a_{14}^{''})^{(1)}(T_{14},t)}$, $\boxed{(a_{15}^{''})^{(1)}(T_{14},t)}$ are first augmentation coefficients for | |

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| category 1, 2 and 3<br>$+(a_{16}'')^{(2,2,)}(T_{17},t)$, $+(a_{17}'')^{(2,2,)}(T_{17},t)$, $+(a_{18}'')^{(2,2,)}(T_{17},t)$ are second augmentation coefficient for category 1, 2 and 3<br>$+(a_{20}'')^{(3,3,)}(T_{21},t)$, $+(a_{21}'')^{(3,3,)}(T_{21},t)$, $+(a_{22}'')^{(3,3,)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3<br>$+(a_{24}'')^{(4,4,4,4,)}(T_{25},t)$, $+(a_{25}'')^{(4,4,4,4,)}(T_{25},t)$, $+(a_{26}'')^{(4,4,4,4,)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3<br>$+(a_{28}'')^{(5,5,5,5,)}(T_{29},t)$, $+(a_{29}'')^{(5,5,5,5,)}(T_{29},t)$, $+(a_{30}'')^{(5,5,5,5,)}(T_{29},t)$ are fifth augmentation coefficient for category 1, 2 and 3<br>$+(a_{32}'')^{(6,6,6,6,)}(T_{33},t)$, $+(a_{33}'')^{(6,6,6,6,)}(T_{33},t)$, $+(a_{34}'')^{(6,6,6,6,)}(T_{33},t)$ are sixth augmentation coefficient for category 1, 2 and 3<br>$+(a_{38}'')^{(7,7)}(T_{37},t)$ $+(a_{37}'')^{(7,7)}(T_{37},t)$ $+(a_{36}'')^{(7,7)}(T_{37},t)$ are seventh augmentation coefficient for 1,2,3<br>$+(a_{40}'')^{(8,8)}(T_{41},t)$ $+(a_{41}'')^{(8,8)}(T_{41},t)$ $+(a_{42}'')^{(8,8)}(T_{41},t)$ are eight augmentation coefficient for 1,2,3<br>$+(a_{44}'')^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{45}'')^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{46}'')^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficient for 1,2,3 | |
| $\frac{dT_{13}}{dt} = (b_{13})^{(1)}T_{14} - \begin{bmatrix} (b_{13}')^{(1)}\left[-(b_{13}'')^{(1)}(G,t)\right] & \left[-(b_{16}'')^{(2,2,)}(G_{19},t)\right] & \left[-(b_{20}'')^{(3,3,)}(G_{23},t)\right] \\ \left[-(b_{24}'')^{(4,4,4,4,)}(G_{27},t)\right] & \left[-(b_{28}'')^{(5,5,5,5,)}(G_{31},t)\right] & \left[-(b_{32}'')^{(6,6,6,6,)}(G_{35},t)\right] \\ \left[-(b_{36}'')^{(7,7,)}(G_{39},t)\right] & \left[-(b_{40}'')^{(8,8)}(G_{43},t)\right] & \left[-(b_{44}'')^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{13}$ | 58 |
| $\frac{dT_{14}}{dt} = (b_{14})^{(1)}T_{13} - \begin{bmatrix} (b_{14}')^{(1)}\left[-(b_{14}'')^{(1)}(G,t)\right] & \left[-(b_{17}'')^{(2,2,)}(G_{19},t)\right] & \left[-(b_{21}'')^{(3,3,)}(G_{23},t)\right] \\ \left[-(b_{25}'')^{(4,4,4,4,)}(G_{27},t)\right] & \left[-(b_{29}'')^{(5,5,5,5,)}(G_{31},t)\right] & \left[-(b_{33}'')^{(6,6,6,6,)}(G_{35},t)\right] \\ \left[-(b_{37}'')^{(7,7,)}(G_{39},t)\right] & \left[-(b_{41}'')^{(8,8)}(G_{43},t)\right] & \left[-(b_{45}'')^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{14}$ | 59 |
| $\frac{dT_{15}}{dt} = (b_{15})^{(1)}T_{14} - \begin{bmatrix} (b_{15}')^{(1)}\left[-(b_{15}'')^{(1)}(G,t)\right] & \left[-(b_{18}'')^{(2,2,)}(G_{19},t)\right] & \left[-(b_{22}'')^{(3,3,)}(G_{23},t)\right] \\ \left[-(b_{26}'')^{(4,4,4,4,)}(G_{27},t)\right] & \left[-(b_{30}'')^{(5,5,5,5,)}(G_{31},t)\right] & \left[-(b_{34}'')^{(6,6,6,6,)}(G_{35},t)\right] \\ \left[-(b_{38}'')^{(7,7,)}(G_{39},t)\right] & \left[-(b_{42}'')^{(8,8)}(G_{43},t)\right] & \left[-(b_{46}'')^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{15}$ | *60* |
| Where $-(b_{13}'')^{(1)}(G,t)$, $-(b_{14}'')^{(1)}(G,t)$, $-(b_{15}'')^{(1)}(G,t)$ are first detrition coefficients for category 1, 2 and 3<br>$-(b_{16}'')^{(2,2,)}(G_{19},t)$, $-(b_{17}'')^{(2,2,)}(G_{19},t)$, $-(b_{18}'')^{(2,2,)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3<br>$-(b_{20}'')^{(3,3,)}(G_{23},t)$, $-(b_{21}'')^{(3,3,)}(G_{23},t)$, $-(b_{22}'')^{(3,3,)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3<br>$-(b_{24}'')^{(4,4,4,4,)}(G_{27},t)$, $-(b_{25}'')^{(4,4,4,4,)}(G_{27},t)$, $-(b_{26}'')^{(4,4,4,4,)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3<br>$-(b_{28}'')^{(5,5,5,5,)}(G_{31},t)$, $-(b_{29}'')^{(5,5,5,5,)}(G_{31},t)$, $-(b_{30}'')^{(5,5,5,5,)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3<br>$-(b_{32}'')^{(6,6,6,6,)}(G_{35},t)$, $-(b_{33}'')^{(6,6,6,6,)}(G_{35},t)$, $-(b_{34}'')^{(6,6,6,6,)}(G_{35},t)$ are sixth detrition coefficients for | |

| | |
|---|---|
| category 1, 2 and 3 $\boxed{-(b_{37}^{''})^{(7,7,)}(G_{39},t)}$, $\boxed{-(b_{36}^{''})^{(7,7,)}(G_{39},t)}$, $\boxed{-(b_{38}^{''})^{(7,7,)}(G_{39},t)}$ are seventh detrition coefficients for category 1, 2 and 3 $\boxed{-(b_{40}^{''})^{(8,8)}(G_{43},t)}$ $\boxed{-(b_{41}^{''})^{(8,8)}(G_{43},t)}$ $\boxed{-(b_{42}^{''})^{(8,8)}(G_{43},t)}$ are eight detrition coefficients for category 1, 2 and 3 $\boxed{-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2 and 3 | |
| | |
| $\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \left[\begin{array}{ccc} (a_{16}^{'})^{(2)} & \boxed{+(a_{16}^{''})^{(2)}(T_{17},t)} & \boxed{+(a_{13}^{''})^{(1,1,)}(T_{14},t)} \quad \boxed{+(a_{20}^{''})^{(3,3,3)}(T_{21},t)} \\ \boxed{+(a_{24}^{''})^{(4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{28}^{''})^{(5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{32}^{''})^{(6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{36}^{''})^{(7,7,7)}(T_{37},t)} & \boxed{+(a_{40}^{''})^{(8,8,8)}(T_{41},t)} & \boxed{+(a_{44}^{''})^{(9,9)}(T_{45},t)} \end{array}\right] G_{16}$ | *61* |
| $\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \left[\begin{array}{ccc} (a_{17}^{'})^{(2)} & \boxed{+(a_{17}^{''})^{(2)}(T_{17},t)} & \boxed{+(a_{14}^{''})^{(1,1,)}(T_{14},t)} \quad \boxed{+(a_{21}^{''})^{(3,3,3)}(T_{21},t)} \\ \boxed{+(a_{25}^{''})^{(4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{29}^{''})^{(5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{33}^{''})^{(6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{37}^{''})^{(7,7,7)}(T_{37},t)} & \boxed{+(a_{41}^{''})^{(8,8,8)}(T_{41},t)} & \boxed{+(a_{45}^{''})^{(9,9)}(T_{45},t)} \end{array}\right] G_{17}$ | *62* |
| $\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \left[\begin{array}{ccc} (a_{18}^{'})^{(2)} & \boxed{+(a_{18}^{''})^{(2)}(T_{17},t)} & \boxed{+(a_{15}^{''})^{(1,1,)}(T_{14},t)} \quad \boxed{+(a_{22}^{''})^{(3,3,3)}(T_{21},t)} \\ \boxed{+(a_{26}^{''})^{(4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{30}^{''})^{(5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{34}^{''})^{(6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{38}^{''})^{(7,7,7)}(T_{37},t)} & \boxed{+(a_{42}^{''})^{(8,8,8)}(T_{41},t)} & \boxed{+(a_{46}^{''})^{(9,9)}(T_{45},t)} \end{array}\right] G_{18}$ | *63* |
| Where $\boxed{+(a_{16}^{''})^{(2)}(T_{17},t)}$, $\boxed{+(a_{17}^{''})^{(2)}(T_{17},t)}$, $\boxed{+(a_{18}^{''})^{(2)}(T_{17},t)}$ are first augmentation coefficients for category 1, 2 and 3 $\boxed{+(a_{13}^{''})^{(1,1,)}(T_{14},t)}$, $\boxed{+(a_{14}^{''})^{(1,1,)}(T_{14},t)}$, $\boxed{+(a_{15}^{''})^{(1,1,)}(T_{14},t)}$ are second augmentation coefficient for category 1, 2 and 3 $\boxed{+(a_{20}^{''})^{(3,3,3)}(T_{21},t)}$, $\boxed{+(a_{21}^{''})^{(3,3,3)}(T_{21},t)}$, $\boxed{+(a_{22}^{''})^{(3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3 $\boxed{+(a_{24}^{''})^{(4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{25}^{''})^{(4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a_{26}^{''})^{(4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3 $\boxed{+(a_{28}^{''})^{(5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{29}^{''})^{(5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a_{30}^{''})^{(5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3 $\boxed{+(a_{32}^{''})^{(6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{33}^{''})^{(6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a_{34}^{''})^{(6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3 $\boxed{+(a_{36}^{''})^{(7,7,7)}(T_{37},t)}$, $\boxed{+(a_{37}^{''})^{(7,7,7)}(T_{37},t)}$, $\boxed{+(a_{38}^{''})^{(7,7,7)}(T_{37},t)}$ are seventh augmentation coefficient for category 1, 2 and 3 $\boxed{+(a_{40}^{''})^{(8,8,8)}(T_{41},t)}$, $\boxed{+(a_{41}^{''})^{(8,8,8)}(T_{41},t)}$, $\boxed{+(a_{42}^{''})^{(8,8,8)}(T_{41},t)}$ are eight augmentation coefficient for category 1, 2 and 3 $\boxed{+(a_{44}^{''})^{(9,9)}(T_{45},t)}$, $\boxed{+(a_{45}^{''})^{(9,9)}(T_{45},t)}$, $\boxed{+(a_{46}^{''})^{(9,9)}(T_{45},t)}$ are ninth augmentation coefficient for category 1, 2 and 3 | |

$$\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \begin{bmatrix} (b'_{16})^{(2)}\boxed{-(b''_{16})^{(2)}(G_{19},t)} & \boxed{-(b''_{13})^{(1,1,)}(G,t)} & \boxed{-(b''_{20})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b''_{24})^{(4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{28})^{(5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{32})^{(6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{36})^{(7,7,7)}(G_{39},t)} & \boxed{-(b''_{40})^{(8,8,8)}(G_{43},t)} & \boxed{-(b''_{44})^{(9,9)}(G_{47},t)} \end{bmatrix} T_{16} \quad (64)$$

$$\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \begin{bmatrix} (b'_{17})^{(2)}\boxed{-(b''_{17})^{(2)}(G_{19},t)} & \boxed{-(b''_{14})^{(1,1,)}(G,t)} & \boxed{-(b''_{21})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b''_{25})^{(4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{29})^{(5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{33})^{(6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{37})^{(7,7,7)}(G_{39},t)} & \boxed{-(b''_{41})^{(8,8,8)}(G_{43},t)} & \boxed{-(b''_{45})^{(9,9)}(G_{47},t)} \end{bmatrix} T_{17} \quad (65)$$

$$\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \begin{bmatrix} (b'_{18})^{(2)}\boxed{-(b''_{18})^{(2)}(G_{19},t)} & \boxed{-(b''_{15})^{(1,1,)}(G,t)} & \boxed{-(b''_{22})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b''_{26})^{(4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{38})^{(7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9)}(G_{47},t)} \end{bmatrix} T_{18} \quad (66)$$

where $\boxed{-(b''_{16})^{(2)}(G_{19},t)}$ , $\boxed{-(b''_{17})^{(2)}(G_{19},t)}$ , $\boxed{-(b''_{18})^{(2)}(G_{19},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{13})^{(1,1,)}(G,t)}$, $\boxed{-(b''_{14})^{(1,1,)}(G,t)}$, $\boxed{-(b''_{15})^{(1,1,)}(G,t)}$ are second detrition coefficients for category 1,2 and 3

$\boxed{-(b''_{20})^{(3,3,3,)}(G_{23},t)}$, $\boxed{-(b''_{21})^{(3,3,3,)}(G_{23},t)}$, $\boxed{-(b''_{22})^{(3,3,3,)}(G_{23},t)}$ are third detrition coefficients for category 1,2 and 3

$\boxed{-(b''_{24})^{(4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{25})^{(4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b''_{26})^{(4,4,4,4,4,4)}(G_{27},t)}$ are fourth detrition coefficients for category 1,2 and 3

$\boxed{-(b''_{28})^{(5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{29})^{(5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b''_{30})^{(5,5,5,5,5,5)}(G_{31},t)}$ are fifth detrition coefficients for category 1,2 and 3

$\boxed{-(b''_{32})^{(6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b''_{33})^{(6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b''_{34})^{(6,6,6,6,6,6)}(G_{35},t)}$ are sixth detrition coefficients for category 1,2 and 3

$\boxed{-(b''_{36})^{(7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{37})^{(7,7,7)}(G_{39},t)}$, $\boxed{-(b''_{38})^{(7,7,7)}(G_{39},t)}$ are seventh detrition coefficients for category 1,2 and 3

$\boxed{-(b''_{40})^{(8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{41})^{(8,8,8)}(G_{43},t)}$, $\boxed{-(b''_{42})^{(8,8,8)}(G_{43},t)}$ are eight detrition coefficients for category 1,2 and 3

$\boxed{-(b''_{44})^{(9,9)}(G_{47},t)}$, $\boxed{-(b''_{46})^{(9,9)}(G_{47},t)}$, $\boxed{-(b''_{45})^{(9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1,2 and 3

$$\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \begin{bmatrix} (a'_{20})^{(3)}\boxed{+(a''_{20})^{(3)}(T_{21},t)} & \boxed{+(a''_{16})^{(2,2,2)}(T_{17},t)} & \boxed{+(a''_{13})^{(1,1,1,)}(T_{14},t)} \\ \boxed{+(a''_{24})^{(4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{36})^{(7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{40})^{(8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{44})^{(9,9,9)}(T_{45},t)} \end{bmatrix} G_{20} \quad (67)$$

$$\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \begin{bmatrix} (a'_{21})^{(3)}\boxed{+(a''_{21})^{(3)}(T_{21},t)} & \boxed{+(a''_{17})^{(2,2,2)}(T_{17},t)} & \boxed{+(a''_{14})^{(1,1,1,)}(T_{14},t)} \\ \boxed{+(a''_{25})^{(4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{29})^{(5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{33})^{(6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{37})^{(7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{41})^{(8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{45})^{(9,9,9)}(T_{45},t)} \end{bmatrix} G_{21} \quad (68)$$

$$\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \begin{bmatrix} (a'_{22})^{(3)}\boxed{+(a''_{22})^{(3)}(T_{21},t)} & \boxed{+(a''_{18})^{(2,2,2)}(T_{17},t)} & \boxed{+(a''_{15})^{(1,1,1,)}(T_{14},t)} \\ \boxed{+(a''_{26})^{(4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{30})^{(5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{34})^{(6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{38})^{(7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{42})^{(8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{46})^{(9,9,9)}(T_{45},t)} \end{bmatrix} G_{22} \quad (69)$$

$\boxed{+(a''_{20})^{(3)}(T_{21},t)}, \boxed{+(a''_{21})^{(3)}(T_{21},t)}, \boxed{+(a''_{22})^{(3)}(T_{21},t)}$ are first augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{16})^{(2,2,2)}(T_{17},t)}, \boxed{+(a''_{17})^{(2,2,2)}(T_{17},t)}, \boxed{+(a''_{18})^{(2,2,2)}(T_{17},t)}$ are second augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{13})^{(1,1,1,)}(T_{14},t)}, \boxed{+(a''_{14})^{(1,1,1,)}(T_{14},t)}, \boxed{+(a''_{15})^{(1,1,1,)}(T_{14},t)}$ are third augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{24})^{(4,4,4,4,4,4)}(T_{25},t)}, \boxed{+(a''_{25})^{(4,4,4,4,4,4)}(T_{25},t)}, \boxed{+(a''_{26})^{(4,4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{28})^{(5,5,5,5,5,5)}(T_{29},t)}, \boxed{+(a''_{29})^{(5,5,5,5,5,5)}(T_{29},t)}, \boxed{+(a''_{30})^{(5,5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{32})^{(6,6,6,6,6,6)}(T_{33},t)}, \boxed{+(a''_{33})^{(6,6,6,6,6,6)}(T_{33},t)}, \boxed{+(a''_{34})^{(6,6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{36})^{(7,7,7,7)}(T_{37},t)}, \boxed{+(a''_{37})^{(7,7,7,7)}(T_{37},t)}, \boxed{+(a''_{38})^{(7,7,7,7)}(T_{37},t)}$ are seventh augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{40})^{(8,8,8,8)}(T_{41},t)}, \boxed{+(a''_{41})^{(8,8,8,8)}(T_{41},t)}, \boxed{+(a''_{42})^{(8,8,8,8)}(T_{41},t)}$ are eight augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{44})^{(9,9,9)}(T_{45},t)}, \boxed{+(a''_{45})^{(9,9,9)}(T_{45},t)}, \boxed{+(a''_{46})^{(9,9,9)}(T_{45},t)}$ are ninth augmentation coefficients for category 1, 2 and 3

$$\frac{dT_{20}}{dt} = (b_{20})^{(3)}T_{21} - \begin{bmatrix} (b'_{20})^{(3)}\boxed{-(b''_{20})^{(3)}(G_{23},t)} & \boxed{-(b''_{16})^{(2,2,2)}(G_{19},t)} & \boxed{-(b''_{13})^{(1,1,1,)}(G,t)} \\ \boxed{-(b''_{24})^{(4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{28})^{(5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{32})^{(6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{36})^{(7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{40})^{(8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{44})^{(9,9,9)}(G_{47},t)} \end{bmatrix} T_{20} \quad (70)$$

$$\frac{dT_{21}}{dt} = (b_{21})^{(3)}T_{20} - \begin{bmatrix} (b'_{21})^{(3)}\boxed{-(b''_{21})^{(3)}(G_{23},t)} & \boxed{-(b''_{17})^{(2,2,2)}(G_{19},t)} & \boxed{-(b''_{14})^{(1,1,1,)}(G,t)} \\ \boxed{-(b''_{25})^{(4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{29})^{(5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{33})^{(6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{37})^{(7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{41})^{(8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{45})^{(9,9,9)}(G_{47},t)} \end{bmatrix} T_{21} \quad (71)$$

$$\frac{dT_{22}}{dt} = (b_{22})^{(3)}T_{21} - \begin{bmatrix} (b'_{22})^{(3)}\boxed{-(b''_{22})^{(3)}(G_{23},t)} & \boxed{-(b''_{18})^{(2,2,2)}(G_{19},t)} & \boxed{-(b''_{15})^{(1,1,1,)}(G,t)} \\ \boxed{-(b''_{26})^{(4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{38})^{(7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9,9)}(G_{47},t)} \end{bmatrix} T_{22} \quad (72)$$

$\boxed{-(b''_{20})^{(3)}(G_{23},t)}, \boxed{-(b''_{21})^{(3)}(G_{23},t)}, \boxed{-(b''_{22})^{(3)}(G_{23},t)}$ *are first detrition coefficients for category 1, 2 and 3*

$\boxed{-(b''_{16})^{(2,2,2)}(G_{19},t)}, \boxed{-(b''_{17})^{(2,2,2)}(G_{19},t)}, \boxed{-(b''_{18})^{(2,2,2)}(G_{19},t)}$ *are second detrition coefficients for category 1, 2 and 3*

$\boxed{-(b''_{13})^{(1,1,1,)}(G,t)}, \boxed{-(b''_{14})^{(1,1,1,)}(G,t)}, \boxed{-(b''_{15})^{(1,1,1,)}(G,t)}$ *are third detrition coefficients for category 1,2 and 3*

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|---|---|
| $-(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$ *are fourth detrition coefficients for category 1, 2 and 3*<br>$-(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$ *are fifth detrition coefficients for category 1, 2 and 3*<br>$-(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$ *are sixth detrition coefficients for category 1, 2 and 3*<br>$-(b_{36}^{''})^{(7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7)}(G_{39},t)$ $-(b_{38}^{''})^{(7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3<br>$-(b_{40}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8)}(G_{43},t)$ are eight detrition coefficients for category 1, 2 and 3<br>$-(b_{46}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3 | |
| $\frac{dG_{24}}{dt} = (a_{24})^{(4)}G_{25} - \begin{bmatrix} (a_{24}^{'})^{(4)} & +(a_{24}^{''})^{(4)}(T_{25},t) & +(a_{28}^{''})^{(5,5,)}(T_{29},t) & +(a_{32}^{''})^{(6,6,)}(T_{33},t) \\ & +(a_{13}^{''})^{(1,1,1,1)}(T_{14},t) & +(a_{16}^{''})^{(2,2,2,2)}(T_{17},t) & +(a_{20}^{''})^{(3,3,3,3)}(T_{21},t) \\ & +(a_{36}^{''})^{(7,7,7,7,7)}(T_{37},t) & +(a_{40}^{''})^{(8,8,8,8,8)}(T_{41},t) & +(a_{44}^{''})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{24}$ | *73* |
| $\frac{dG_{25}}{dt} = (a_{25})^{(4)}G_{24} - \begin{bmatrix} (a_{25}^{'})^{(4)} & +(a_{25}^{''})^{(4)}(T_{25},t) & +(a_{29}^{''})^{(5,5,)}(T_{29},t) & +(a_{33}^{''})^{(6,6)}(T_{33},t) \\ & +(a_{14}^{''})^{(1,1,1,1)}(T_{14},t) & +(a_{17}^{''})^{(2,2,2,2)}(T_{17},t) & +(a_{21}^{''})^{(3,3,3,3)}(T_{21},t) \\ & +(a_{37}^{''})^{(7,7,7,7,7)}(T_{37},t) & +(a_{41}^{''})^{(8,8,8,8,8)}(T_{41},t) & +(a_{45}^{''})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{25}$ | *74* |
| $\frac{dG_{26}}{dt} = (a_{26})^{(4)}G_{25} - \begin{bmatrix} (a_{26}^{'})^{(4)} & +(a_{26}^{''})^{(4)}(T_{25},t) & +(a_{30}^{''})^{(5,5,)}(T_{29},t) & +(a_{34}^{''})^{(6,6,)}(T_{33},t) \\ & +(a_{15}^{''})^{(1,1,1,1)}(T_{14},t) & +(a_{18}^{''})^{(2,2,2,2)}(T_{17},t) & +(a_{22}^{''})^{(3,3,3,3)}(T_{21},t) \\ & +(a_{38}^{''})^{(7,7,7,7,7)}(T_{37},t) & +(a_{42}^{''})^{(8,8,8,8,8)}(T_{41},t) & +(a_{46}^{''})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{26}$ | *75* |

$(a_{24}^{''})^{(4)}(T_{25},t)$, $(a_{25}^{''})^{(4)}(T_{25},t)$, $(a_{26}^{''})^{(4)}(T_{25},t)$ are first augmentation coefficients category 1, 2 3

$+(a_{28}^{''})^{(5,5,)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,)}(T_{29},t)$ are second augmentation coefficient for category 1, 2 and 3

$+(a_{32}^{''})^{(6,6,)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,)}(T_{33},t)$ are third augmentation coefficient for category 1, 2 and 3

$+(a_{13}^{''})^{(1,1,1,1)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1)}(T_{14},t)$ are fourth augmentation coefficients for category 1, 2 and 3

$+(a_{16}^{''})^{(2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2)}(T_{17},t)$ are fifth augmentation coefficients for category 1, 2 and 3

$+(a_{20}^{''})^{(3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3)}(T_{21},t)$ are sixth augmentation coefficients for category 1, 2 and 3

$+(a_{36}^{''})^{(7,7,7,7,7)}(T_{37},t)$, $+(a_{37}^{''})^{(7,7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7,7)}(T_{37},t)$ are seventh augmentation coefficients for category 1, 2 and 3

$+(a_{40}^{''})^{(8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8,8)}(T_{41},t)$ are eighth augmentation coefficients for category 1, 2 and 3

$+(a_{46}^{''})^{(9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9)}(T_{45},t)$ are ninth detrition coefficients for category 1 2 3

$$\frac{dT_{24}}{dt} = (b_{24})^{(4)}T_{25} - \begin{bmatrix} (b_{24}^{'})^{(4)} & -(b_{24}^{''})^{(4)}(G_{27},t) & -(b_{28}^{''})^{(5,5,)}(G_{31},t) & -(b_{32}^{''})^{(6,6,)}(G_{35},t) \\ & -(b_{13}^{''})^{(1,1,1,1)}(G,t) & -(b_{16}^{''})^{(2,2,2,2)}(G_{19},t) & -(b_{20}^{''})^{(3,3,3,3)}(G_{23},t) \\ & -(b_{36}^{''})^{(7,7,7,7,7)}(G_{39},t) & -(b_{40}^{''})^{(8,8,8,8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{24} \quad (76)$$

$$\frac{dT_{25}}{dt} = (b_{25})^{(4)}T_{24} - \begin{bmatrix} (b_{25}^{'})^{(4)} & -(b_{25}^{''})^{(4)}(G_{27},t) & -(b_{29}^{''})^{(5,5,)}(G_{31},t) & -(b_{33}^{''})^{(6,6,)}(G_{35},t) \\ & -(b_{14}^{''})^{(1,1,1,1)}(G,t) & -(b_{17}^{''})^{(2,2,2,2)}(G_{19},t) & -(b_{21}^{''})^{(3,3,3,3)}(G_{23},t) \\ & -(b_{37}^{''})^{(7,7,7,7,7)}(G_{39},t) & -(b_{41}^{''})^{(8,8,8,8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{25} \quad (77)$$

$$\frac{dT_{26}}{dt} = (b_{26})^{(4)}T_{25} - \begin{bmatrix} (b_{26}^{'})^{(4)} & -(b_{26}^{''})^{(4)}(G_{27},t) & -(b_{30}^{''})^{(5,5,)}(G_{31},t) & -(b_{34}^{''})^{(6,6,)}(G_{35},t) \\ & -(b_{15}^{''})^{(1,1,1,1)}(G,t) & -(b_{18}^{''})^{(2,2,2,2)}(G_{19},t) & -(b_{22}^{''})^{(3,3,3,3)}(G_{23},t) \\ & -(b_{38}^{''})^{(7,7,7,7,7)}(G_{39},t) & -(b_{42}^{''})^{(8,8,8,8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{26} \quad (78)$$

Where $-(b_{24}^{''})^{(4)}(G_{27},t)$, $-(b_{25}^{''})^{(4)}(G_{27},t)$, $-(b_{26}^{''})^{(4)}(G_{27},t)$ are first detrition coefficients for category 1, 2 and 3

$-(b_{28}^{''})^{(5,5,)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,)}(G_{31},t)$ are second detrition coefficients for category 1, 2 and 3

$-(b_{32}^{''})^{(6,6,)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,)}(G_{35},t)$ are third detrition coefficients for category 1, 2 and 3

$-(b_{13}^{''})^{(1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1)}(G,t)$, $-(b_{15}^{''})^{(1,1,1,1)}(G,t)$ are fourth detrition coefficients for category 1, 2 and 3

$-(b_{16}^{''})^{(2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2)}(G_{19},t)$ are fifth detrition coefficients for category 1, 2 and 3

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|---|---|
| $-(b_{20}^{''})^{(3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3)}(G_{23},t)$ *are sixth detrition coefficients for category* 1, 2 *and* 3 $-(b_{36}^{''})^{(7,7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7,7)}(G_{39},t)$ *are seventh detrition coefficients for category* 1, 2 *and* 3 $-(b_{40}^{''})^{(8,8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8,8)}(G_{43},t)$ *are eighth detrition coefficients for category* 1, 2 *and* 3 $-(b_{46}^{''})^{(9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1 2 3 | |
| $\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \begin{bmatrix} (a_{28}^{'})^{(5)} & +(a_{28}^{''})^{(5)}(T_{29},t) & +(a_{24}^{''})^{(4,4,)}(T_{25},t) & +(a_{32}^{''})^{(6,6,6)}(T_{33},t) \\ & +(a_{13}^{''})^{(1,1,1,1,1)}(T_{14},t) & +(a_{16}^{''})^{(2,2,2,2,2)}(T_{17},t) & +(a_{20}^{''})^{(3,3,3,3,3)}(T_{21},t) \\ & +(a_{36}^{''})^{(7,7,7,7,7,7)}(T_{37},t) & +(a_{40}^{''})^{(8,8,8,8,8,8)}(T_{41},t) & +(a_{44}^{''})^{(9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{28}$ | *79* |
| $\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \begin{bmatrix} (a_{29}^{'})^{(5)} & +(a_{29}^{''})^{(5)}(T_{29},t) & +(a_{25}^{''})^{(4,4,)}(T_{25},t) & +(a_{33}^{''})^{(6,6,6)}(T_{33},t) \\ & +(a_{14}^{''})^{(1,1,1,1,1)}(T_{14},t) & +(a_{17}^{''})^{(2,2,2,2,2)}(T_{17},t) & +(a_{21}^{''})^{(3,3,3,3,3)}(T_{21},t) \\ & +(a_{37}^{''})^{(7,7,7,7,7,7)}(T_{37},t) & +(a_{41}^{''})^{(8,8,8,8,8,8)}(T_{41},t) & +(a_{45}^{''})^{(9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{29}$ | *80* |
| $\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \begin{bmatrix} (a_{30}^{'})^{(5)} & +(a_{30}^{''})^{(5)}(T_{29},t) & +(a_{26}^{''})^{(4,4,)}(T_{25},t) & +(a_{34}^{''})^{(6,6,6)}(T_{33},t) \\ & +(a_{15}^{''})^{(1,1,1,1,1)}(T_{14},t) & +(a_{18}^{''})^{(2,2,2,2,2)}(T_{17},t) & +(a_{22}^{''})^{(3,3,3,3,3)}(T_{21},t) \\ & +(a_{38}^{''})^{(7,7,7,7,7,7)}(T_{37},t) & +(a_{42}^{''})^{(8,8,8,8,8,8)}(T_{41},t) & +(a_{46}^{''})^{(9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{30}$ | *81* |
| *Where* $+(a_{28}^{''})^{(5)}(T_{29},t)$, $+(a_{29}^{''})^{(5)}(T_{29},t)$, $+(a_{30}^{''})^{(5)}(T_{29},t)$ *are first augmentation coefficients for category* 1, 2 *and* 3 *And* $+(a_{24}^{''})^{(4,4,)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,)}(T_{25},t)$ *are second augmentation coefficient for category* 1, 2 *and* 3 $+(a_{32}^{''})^{(6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6)}(T_{33},t)$ *are third augmentation coefficient for category* 1, 2 *and* 3 $+(a_{13}^{''})^{(1,1,1,1,1)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1,1)}(T_{14},t)$ are fourth augmentation coefficients for category 1,2, and 3 $+(a_{16}^{''})^{(2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2)}(T_{17},t)$ are fifth augmentation coefficients for category 1,2,and 3 $+(a_{20}^{''})^{(3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3)}(T_{21},t)$ are sixth augmentation coefficients for category 1,2, 3 $+(a_{36}^{''})^{(7,7,7,7,7,7)}(T_{37},t)$, $+(a_{37}^{''})^{(7,7,7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7,7,7)}(T_{37},t)$ are seventh augmentation coefficients for category 1,2, 3 $+(a_{40}^{''})^{(8,8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8,8,8)}(T_{41},t)$ are eighth augmentation coefficients for category 1,2, 3 $+(a_{46}^{''})^{(9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficients for category 1,2, 3 | |

$$\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \begin{bmatrix} (b'_{28})^{(5)} & -(b''_{28})^{(5)}(G_{31},t) & -(b''_{24})^{(4,4,)}(G_{27},t) & -(b''_{32})^{(6,6,6)}(G_{35},t) \\ -(b''_{13})^{(1,1,1,1,1)}(G,t) & -(b''_{16})^{(2,2,2,2,2)}(G_{19},t) & -(b''_{20})^{(3,3,3,3,3)}(G_{23},t) & \\ -(b''_{36})^{(7,7,7,7,7,7)}(G_{39},t) & -(b''_{40})^{(8,8,8,8,8,8)}(G_{43},t) & -(b''_{44})^{(9,9,9,9,9)}(G_{47},t) & \end{bmatrix} T_{28} \quad (82)$$

$$\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \begin{bmatrix} (b'_{29})^{(5)} & -(b''_{29})^{(5)}(G_{31},t) & -(b''_{25})^{(4,4,)}(G_{27},t) & -(b''_{33})^{(6,6,6)}(G_{35},t) \\ -(b''_{14})^{(1,1,1,1,1)}(G,t) & -(b''_{17})^{(2,2,2,2,2)}(G_{19},t) & -(b''_{21})^{(3,3,3,3,3)}(G_{23},t) & \\ -(b''_{37})^{(7,7,7,7,7,7)}(G_{39},t) & -(b''_{41})^{(8,8,8,8,8,8)}(G_{43},t) & -(b''_{45})^{(9,9,9,9,9)}(G_{47},t) & \end{bmatrix} T_{29} \quad (83)$$

$$\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \begin{bmatrix} (b'_{30})^{(5)} & -(b''_{30})^{(5)}(G_{31},t) & -(b''_{26})^{(4,4,)}(G_{27},t) & -(b''_{34})^{(6,6,6)}(G_{35},t) \\ -(b''_{15})^{(1,1,1,1,1,)}(G,t) & -(b''_{18})^{(2,2,2,2,2)}(G_{19},t) & -(b''_{22})^{(3,3,3,3,3)}(G_{23},t) & \\ -(b''_{38})^{(7,7,7,7,7,7)}(G_{39},t) & -(b''_{42})^{(8,8,8,8,8,8)}(G_{43},t) & -(b''_{46})^{(9,9,9,9,9)}(G_{47},t) & \end{bmatrix} T_{30} \quad (84)$$

*where* $-(b''_{28})^{(5)}(G_{31},t)$ , $-(b''_{29})^{(5)}(G_{31},t)$ , $-(b''_{30})^{(5)}(G_{31},t)$ *are first detrition coefficients for category* 1, 2 *and* 3

$-(b''_{24})^{(4,4,)}(G_{27},t)$, $-(b''_{25})^{(4,4,)}(G_{27},t)$, $-(b''_{26})^{(4,4,)}(G_{27},t)$ *are second detrition coefficients for category* 1,2 *and* 3

$-(b''_{32})^{(6,6,6)}(G_{35},t)$, $-(b''_{33})^{(6,6,6)}(G_{35},t)$, $-(b''_{34})^{(6,6,6)}(G_{35},t)$ *are third detrition coefficients for category* 1,2 *and* 3

$-(b''_{13})^{(1,1,1,1,1)}(G,t)$, $-(b''_{14})^{(1,1,1,1,1)}(G,t)$ , $-(b''_{15})^{(1,1,1,1,1,)}(G,t)$ are fourth detrition coefficients for category 1,2, and 3

$-(b''_{16})^{(2,2,2,2,2)}(G_{19},t)$, $-(b''_{17})^{(2,2,2,2,2)}(G_{19},t)$, $-(b''_{18})^{(2,2,2,2,2)}(G_{19},t)$ are fifth detrition coefficients for category 1,2, and 3

$-(b''_{20})^{(3,3,3,3,3)}(G_{23},t)$, $-(b''_{21})^{(3,3,3,3,3)}(G_{23},t)$, $-(b''_{22})^{(3,3,3,3,3)}(G_{23},t)$ are sixth detrition coefficients for category 1,2, and 3

$-(b''_{36})^{(7,7,7,7,7,7)}(G_{39},t)$, $-(b''_{37})^{(7,7,7,7,7,7)}(G_{39},t)$, $-(b''_{38})^{(7,7,7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1,2, and 3

$-(b''_{42})^{(8,8,8,8,8,8)}(G_{43},t)$, $-(b''_{41})^{(8,8,8,8,8,8)}(G_{43},t)$, $-(b''_{40})^{(8,8,8,8,8,8)}(G_{43},t)$ are eighth detrition coefficients for category 1,2, and 3

$-(b''_{46})^{(9,9,9,9,9)}(G_{47},t)$, $-(b''_{45})^{(9,9,9,9,9)}(G_{47},t)$, $-(b''_{44})^{(9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1,2, and 3

$$\frac{dG_{32}}{dt} = (a_{32})^{(6)}G_{33} - \begin{bmatrix} (a'_{32})^{(6)} & +(a''_{32})^{(6)}(T_{33},t) & +(a''_{28})^{(5,5,5)}(T_{29},t) & +(a''_{24})^{(4,4,4,)}(T_{25},t) \\ +(a''_{13})^{(1,1,1,1,1,1)}(T_{14},t) & +(a''_{16})^{(2,2,2,2,2,2)}(T_{17},t) & +(a''_{20})^{(3,3,3,3,3,3)}(T_{21},t) & \\ +(a''_{36})^{(7,7,7,7,7,7,7)}(T_{37},t) & +(a''_{40})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a''_{44})^{(9,9,9,9,9,9)}(T_{45},t) & \end{bmatrix} G_{32} \quad (85)$$

$$\frac{dG_{33}}{dt} = (a_{33})^{(6)}G_{32} - \begin{bmatrix} (a'_{33})^{(6)} & +(a''_{33})^{(6)}(T_{33},t) & +(a''_{29})^{(5,5,5)}(T_{29},t) & +(a''_{25})^{(4,4,4,)}(T_{25},t) \\ +(a''_{14})^{(1,1,1,1,1,1)}(T_{14},t) & +(a''_{17})^{(2,2,2,2,2,2)}(T_{17},t) & +(a''_{21})^{(3,3,3,3,3,3)}(T_{21},t) & \\ +(a''_{37})^{(7,7,7,7,7,7,7)}(T_{37},t) & +(a''_{41})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a''_{45})^{(9,9,9,9,9,9)}(T_{45},t) & \end{bmatrix} G_{33} \quad (86)$$

$$\frac{dG_{34}}{dt} = (a_{34})^{(6)} G_{33} - \begin{bmatrix} (a'_{34})^{(6)} & +(a''_{34})^{(6)}(T_{33},t) & +(a''_{30})^{(5,5,5)}(T_{29},t) & +(a''_{26})^{(4,4,4,)}(T_{25},t) \\ & +(a''_{15})^{(1,1,1,1,1,1)}(T_{14},t) & +(a''_{18})^{(2,2,2,2,2,2)}(T_{17},t) & +(a''_{22})^{(3,3,3,3,3,3)}(T_{21},t) \\ & +(a''_{38})^{(7,7,7,7,7,7,7)}(T_{37},t) & +(a''_{42})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a''_{46})^{(9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{34} \quad (87)$$

$+(a''_{32})^{(6)}(T_{33},t)$, $+(a''_{33})^{(6)}(T_{33},t)$, $+(a''_{34})^{(6)}(T_{33},t)$ *are first augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{28})^{(5,5,5)}(T_{29},t)$, $+(a''_{29})^{(5,5,5)}(T_{29},t)$, $+(a''_{30})^{(5,5,5)}(T_{29},t)$ *are second augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{24})^{(4,4,4,)}(T_{25},t)$, $+(a''_{25})^{(4,4,4,)}(T_{25},t)$, $+(a''_{26})^{(4,4,4,)}(T_{25},t)$ *are third augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{13})^{(1,1,1,1,1,1)}(T_{14},t)$, $+(a''_{14})^{(1,1,1,1,1,1)}(T_{14},t)$, $+(a''_{15})^{(1,1,1,1,1,1)}(T_{14},t)$ - are fourth augmentation coefficients

$+(a''_{16})^{(2,2,2,2,2,2)}(T_{17},t)$, $+(a''_{17})^{(2,2,2,2,2,2)}(T_{17},t)$, $+(a''_{18})^{(2,2,2,2,2,2)}(T_{17},t)$ - fifth augmentation coefficients

$+(a''_{20})^{(3,3,3,3,3,3)}(T_{21},t)$, $+(a''_{21})^{(3,3,3,3,3,3)}(T_{21},t)$, $+(a''_{22})^{(3,3,3,3,3,3)}(T_{21},t)$ sixth augmentation coefficients

$+(a''_{36})^{(7,7,7,7,7,7,7)}(T_{37},t)$, $+(a''_{37})^{(7,7,7,7,7,7,7)}(T_{37},t)$, $+(a''_{38})^{(7,7,7,7,7,7,7)}(T_{37},t)$ seventh augmentation coefficients

$+(a''_{40})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a''_{41})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a''_{42})^{(8,8,8,8,8,8,8)}(T_{41},t)$ Eighth augmentation coefficients

$+(a''_{44})^{(9,9,9,9,9,9)}(T_{45},t)$, $+(a''_{45})^{(9,9,9,9,9,9)}(T_{45},t)$, $+(a''_{46})^{(9,9,9,9,9,9)}(T_{45},t)$ ninth augmentation coefficients

$$\frac{dT_{32}}{dt} = (b_{32})^{(6)} T_{33} - \begin{bmatrix} (b'_{32})^{(6)} & -(b''_{32})^{(6)}(G_{35},t) & -(b''_{28})^{(5,5,5)}(G_{31},t) & -(b''_{24})^{(4,4,4,)}(G_{27},t) \\ & -(b''_{13})^{(1,1,1,1,1,1)}(G,t) & -(b''_{16})^{(2,2,2,2,2,2)}(G_{19},t) & -(b''_{20})^{(3,3,3,3,3,3)}(G_{23},t) \\ & -(b''_{36})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b''_{40})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b''_{44})^{(9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{32} \quad (88)$$

$$\frac{dT_{33}}{dt} = (b_{33})^{(6)} T_{32} - \begin{bmatrix} (b'_{33})^{(6)} & -(b''_{33})^{(6)}(G_{35},t) & -(b''_{29})^{(5,5,5)}(G_{31},t) & -(b''_{25})^{(4,4,4,)}(G_{27},t) \\ & -(b''_{14})^{(1,1,1,1,1,1)}(G,t) & -(b''_{17})^{(2,2,2,2,2,2)}(G_{19},t) & -(b''_{21})^{(3,3,3,3,3,3)}(G_{23},t) \\ & -(b''_{37})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b''_{41})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b''_{45})^{(9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{33} \quad (89)$$

$$\frac{dT_{34}}{dt} = (b_{34})^{(6)} T_{33} - \begin{bmatrix} (b'_{34})^{(6)} & -(b''_{34})^{(6)}(G_{35},t) & -(b''_{30})^{(5,5,5)}(G_{31},t) & -(b''_{26})^{(4,4,4,)}(G_{27},t) \\ & -(b''_{15})^{(1,1,1,1,1,1)}(G,t) & -(b''_{18})^{(2,2,2,2,2,2)}(G_{19},t) & -(b''_{22})^{(3,3,3,3,3,3)}(G_{23},t) \\ & -(b''_{38})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b''_{42})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b''_{46})^{(9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{34} \quad (90)$$

$-(b''_{32})^{(6)}(G_{35},t)$, $-(b''_{33})^{(6)}(G_{35},t)$, $-(b''_{34})^{(6)}(G_{35},t)$ *are first detrition coefficients for category* 1, 2 *and* 3

$-(b''_{28})^{(5,5,5)}(G_{31},t)$, $-(b''_{29})^{(5,5,5)}(G_{31},t)$, $-(b''_{30})^{(5,5,5)}(G_{31},t)$ *are second detrition coefficients for category* 1, 2 *and* 3

$-(b_{24}^{''})^{(4,4,4,)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,)}(G_{27},t)$ *are third detrition coefficients for category* 1,2 *and* 3

$-(b_{13}^{''})^{(1,1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1,1)}(G,t)$, $-(b_{15}^{''})^{(1,1,1,1,1,1)}(G,t)$ are fourth detrition coefficients for category 1, 2, and 3

$-(b_{16}^{''})^{(2,2,2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2,2,2)}(G_{19},t)$ are fifth detrition coefficients for category 1, 2, and 3

$-(b_{20}^{''})^{(3,3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3,3)}(G_{23},t)$ are sixth detrition coefficients for category 1, 2, and 3

$-(b_{36}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2, and 3

$-(b_{40}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t)$ are eighth detrition coefficients for category 1, 2, and 3

$-(b_{46}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2, and 3

$$\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \begin{bmatrix} (a_{36}^{'})^{(7)}+(a_{36}^{''})^{(7)}(T_{37},t) & +(a_{16}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t) & +(a_{20}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{24}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t) & +(a_{28}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t) & +(a_{32}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{13}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t) & +(a_{40}^{''})^{(8,8,8,8,8,8,8,8,)}(T_{41},t) & +(a_{44}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{13} \quad (91)$$

$$\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \begin{bmatrix} (a_{37}^{'})^{(7)}+(a_{37}^{''})^{(7)}(T_{37},t) & +(a_{17}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t) & +(a_{21}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{25}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t) & +(a_{29}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t) & +(a_{33}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{13}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t) & +(a_{41}^{''})^{(8,8,8,8,8,8,8,8)}(T_{41},t) & +(a_{45}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{14} \quad (92)$$

$$\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \begin{bmatrix} (a_{38}^{'})^{(7)}+(a_{38}^{''})^{(7)}(T_{37},t) & +(a_{18}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t) & +(a_{22}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{26}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t) & +(a_{30}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t) & +(a_{34}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{15}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t) & +(a_{42}^{''})^{(8,8,8,8,8,8,8,8)}(T_{41},t) & +(a_{46}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{15} \quad (93)$$

Where $(a_{36}^{''})^{(7)}(T_{37},t)$, $(a_{37}^{''})^{(7)}(T_{37},t)$, $(a_{38}^{''})^{(7)}(T_{37},t)$ are first augmentation coefficients for category 1, 2 and 3

$+(a_{16}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t)$ are second augmentation coefficient for category 1, 2 and 3

$+(a_{20}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3

$+(a_{24}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3

$+(a_{28}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation

coefficient for category 1, 2 and 3 $+(a_{32}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation coefficient for category 1, 2 and 3 $+(a_{13}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{13}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t)$ are seventh augmentation coefficient for category 1, 2 and 3 $+(a_{42}^{''})^{(8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{40}^{''})^{(8,8,8,8,8,8,8,8,)}(T_{41},t)$ are eighth augmentation coefficient for 1,2,3 $+(a_{46}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{36}}{dt} = (b_{36})^{(7)}T_{37} - \begin{bmatrix} (b_{36}^{'})^{(7)}-(b_{36}^{''})^{(7)}(G_{39},t) & -(b_{16}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t) & -(b_{20}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t) \\ -(b_{24}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t) & -(b_{28}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t) & -(b_{32}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t) \\ -(b_{13}^{''})^{(1,1,1,1,1,1,1)}(G,t) & -(b_{40}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{13}$$ 94

$$\frac{dT_{37}}{dt} = (b_{37})^{(7)}T_{36} - \begin{bmatrix} (b_{37}^{'})^{(7)}-(b_{37}^{''})^{(7)}(G_{39},t) & -(b_{17}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t) & -(b_{21}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t) \\ -(b_{25}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t) & -(b_{29}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t) & -(b_{33}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t) \\ -(b_{14}^{''})^{(1,1,1,1,1,1,1)}(G,t) & -(b_{41}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{14}$$

$$\frac{dT_{38}}{dt} = (b_{38})^{(7)}T_{37} - \begin{bmatrix} (b_{38}^{'})^{(7)}-(b_{38}^{''})^{(7)}(G_{39},t) & -(b_{18}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t) & -(b_{22}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t) \\ -(b_{26}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t) & -(b_{30}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t) & -(b_{34}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t) \\ -(b_{15}^{''})^{(1,1,1,1,1,1,1)}(G,t) & -(b_{42}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{15}$$

Where $-(b_{36}^{''})^{(7)}(G_{39},t)$, $-(b_{37}^{''})^{(7)}(G_{39},t)$, $-(b_{38}^{''})^{(7)}(G_{39},t)$ are first detrition coefficients for category 1, 2 and 3 $-(b_{16}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3 $-(b_{20}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3 $-(b_{24}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3 $-(b_{28}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3 $-(b_{32}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t)$ are sixth detrition coefficients for category 1, 2 and 3 $-(b_{15}^{''})^{(1,1,1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1,1,1)}(G,t)$, $-(b_{13}^{''})^{(1,1,1,1,1,1,1)}(G,t)$ are seventh detrition coefficients for category 1, 2 and 3

$-(b_{40}'')^{(8,8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{41}'')^{(8,8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{42}'')^{(8,8,8,8,8,8,8,8)}(G_{43},t)$ are eighth detrition coefficients for category 1, 2 and 3

$-(b_{46}'')^{(9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}'')^{(9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{44}'')^{(9,9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3

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$$\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \begin{bmatrix} (a_{40}')^{(8)}+(a_{40}'')^{(8)}(T_{41},t) & +(a_{16}'')^{(2,2,2,2,2,2,2,2)}(T_{17},t) & +(a_{20}'')^{(3,3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{24}'')^{(4,4,4,4,4,4,4,4)}(T_{25},t) & +(a_{28}'')^{(5,5,5,5,5,5,5,5)}(T_{29},t) & +(a_{32}'')^{(6,6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{13}'')^{(1,1,1,1,1,1,1,1)}(T_{14},t) & +(a_{36}'')^{(7,7,7,7,7,7,7,7)}(T_{37},t) & +(a_{44}'')^{(9,9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{13}$$

$$\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \begin{bmatrix} (a_{41}')^{(8)}+(a_{41}'')^{(8)}(T_{41},t) & +(a_{17}'')^{(2,2,2,2,2,2,2,2)}(T_{17},t) & +(a_{21}'')^{(3,3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{25}'')^{(4,4,4,4,4,4,4,4)}(T_{25},t) & +(a_{29}'')^{(5,5,5,5,5,5,5,5)}(T_{29},t) & +(a_{33}'')^{(6,6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{13}'')^{(1,1,1,1,1,1,1,1)}(T_{14},t) & +(a_{37}'')^{(7,7,7,7,7,7,7,7)}(T_{37},t) & +(a_{45}'')^{(9,9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{14}$$

$$\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \begin{bmatrix} (a_{42}')^{(8)}+(a_{42}'')^{(8)}(T_{41},t) & +(a_{18}'')^{(2,2,2,2,2,2,2,2)}(T_{17},t) & +(a_{22}'')^{(3,3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{26}'')^{(4,4,4,4,4,4,4,4)}(T_{25},t) & +(a_{30}'')^{(5,5,5,5,5,5,5,5)}(T_{29},t) & +(a_{34}'')^{(6,6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{15}'')^{(1,1,1,1,1,1,1,1)}(T_{14},t) & +(a_{38}'')^{(7,7,7,7,7,7,7,7)}(T_{37},t) & +(a_{46}'')^{(9,9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{15}$$

Where $+(a_{40}'')^{(8)}(T_{41},t)$, $+(a_{41}'')^{(8)}(T_{41},t)$, $+(a_{42}'')^{(8)}(T_{41},t)$ are first augmentation coefficients for category 1, 2 and 3

$+(a_{16}'')^{(2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{17}'')^{(2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{18}'')^{(2,2,2,2,2,2,2,2)}(T_{17},t)$ are second augmentation coefficient for category 1, 2 and 3

$+(a_{20}'')^{(3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{21}'')^{(3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{22}'')^{(3,3,3,3,3,3,3,3)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3

$+(a_{24}'')^{(4,4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}'')^{(4,4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}'')^{(4,4,4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3

$+(a_{28}'')^{(5,5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}'')^{(5,5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}'')^{(5,5,5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficient for category 1, 2 and 3

$+(a_{32}'')^{(6,6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}'')^{(6,6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}'')^{(6,6,6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation coefficient for category 1, 2 and 3

$+(a_{13}'')^{(1,1,1,1,1,1,1,1)}(T_{14},t)$ $+(a_{14}'')^{(1,1,1,1,1,1,1,1)}(T_{14},t)$ $+(a_{15}'')^{(1,1,1,1,1,1,1,1)}(T_{14},t)$ are seventh augmentation coefficient for 1,2,3

$+(a_{36}'')^{(7,7,7,7,7,7,7,7)}(T_{37},t)$, $+(a_{37}'')^{(7,7,7,7,7,7,7,7)}(T_{37},t)$, $+(a_{38}'')^{(7,7,7,7,7,7,7,7)}(T_{37},t)$ are eighth augmentation coefficient for 1,2,3

$+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{40}}{dt} = (b_{40})^{(8)}T_{41} - \begin{bmatrix} (b_{40}^{'})^{(8)}\left[-(b_{40}^{''})^{(8)}(G_{43},t)\right] & -(b_{16}^{''})^{(2,2,2,2,2,2,2,2)}(G_{19},t) & -(b_{20}^{''})^{(3,3,3,3,3,3,3,3)}(G_{23},t) \\ -(b_{24}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t) & -(b_{28}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t) & -(b_{32}^{''})^{(6,6,6,6,6,6,6,6)}(G_{35},t) \\ -(b_{13}^{''})^{(1,1,1,1,1,1,1,1)}(G,t) & -(b_{36}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t) & -(b_{44}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{13}$$

$$\frac{dT_{41}}{dt} = (b_{41})^{(8)}T_{40} - \begin{bmatrix} (b_{41}^{'})^{(8)}\left[-(b_{41}^{''})^{(8)}(G_{43},t)\right] & -(b_{17}^{''})^{(2,2,2,2,2,2,2,2)}(G_{19},t) & -(b_{21}^{''})^{(3,3,3,3,3,3,3,3)}(G_{23},t) \\ -(b_{25}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t) & -(b_{29}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t) & -(b_{33}^{''})^{(6,6,6,6,6,6,6,6)}(G_{35},t) \\ -(b_{14}^{''})^{(1,1,1,1,1,1,1,1)}(G,t) & -(b_{37}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t) & -(b_{45}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{14}$$

$$\frac{dT_{42}}{dt} = (b_{42})^{(8)}T_{41} - \begin{bmatrix} (b_{42}^{'})^{(8)}\left[-(b_{42}^{''})^{(8)}(G_{43},t)\right] & -(b_{18}^{''})^{(2,2,2,2,2,2,2,2)}(G_{19},t) & -(b_{22}^{''})^{(3,3,3,3,3,3,3,3)}(G_{23},t) \\ -(b_{26}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t) & -(b_{30}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t) & -(b_{34}^{''})^{(6,6,6,6,6,6,6,6)}(G_{35},t) \\ -(b_{15}^{''})^{(1,1,1,1,1,1,1,1)}(G,t) & -(b_{38}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t) & -(b_{46}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{15}$$

Where $-(b_{36}^{''})^{(7)}(G_{39},t)$, $-(b_{37}^{''})^{(7)}(G_{39},t)$, $-(b_{38}^{''})^{(7)}(G_{39},t)$ are first detrition coefficients for category 1, 2 and 3

$-(b_{16}^{''})^{(2,2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2,2,2,2,2)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3

$-(b_{20}^{''})^{(3,3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3,3,3,3)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3

$-(b_{24}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3

$-(b_{28}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3

$-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{15}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$ are sixth detrition coefficients for category 1, 2 and 3

$-(b_{13}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{38}^{''})^{(7,7,)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3

$-(b_{36}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)$ are eighth detrition coefficients for category 1, 2 and 3

$-(b_{44}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3



$$\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \left[\begin{array}{ccc} (a'_{44})^{(9)}\boxed{+(a''_{44})^{(9)}(T_{45},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{24})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{36})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)} \end{array}\right] G_{13}$$

$$\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \left[\begin{array}{ccc} (a'_{45})^{(9)}\boxed{+(a''_{45})^{(9)}(T_{45},t)} & \boxed{+(a''_{17})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{21})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{25})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{29})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{33})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{14})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{37})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{41})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)} \end{array}\right] G_{14}$$

$$\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \left[\begin{array}{ccc} (a'_{46})^{(9)}\boxed{+(a''_{46})^{(9)}(T_{37},t)} & \boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)} & \boxed{+(a''_{42})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)} \end{array}\right] G_{15}$$

Where $\boxed{+(a''_{44})^{(9)}(T_{45},t)}$, $\boxed{+(a''_{45})^{(9)}(T_{45},t)}$, $\boxed{+(a''_{46})^{(9)}(T_{37},t)}$ are first augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{16})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{17})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)}$, $\boxed{+(a''_{18})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)}$ are second augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{20})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{21})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{22})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{24})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{25})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{26})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{28})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{29})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{30})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{32})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{33})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{34})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{13})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a''_{14})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)}$, $\boxed{+(a''_{15})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)}$ are Seventh augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{38})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)}$ $\boxed{+(a''_{37})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)}$ $\boxed{+(a''_{36})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)}$ are eighth augmentation coefficient for 1,2,3

$\boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a''_{42})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)}$, $\boxed{+(a''_{41})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)}$ are ninth augmentation coefficient for 1,2,3

$\frac{dT_{44}}{dt} =$

$$(b_{44})^{(9)}T_{45} - \left[\begin{array}{ccc} (b_{44}^{'})^{(9)}\boxed{-(b_{44}^{''})^{(9)}(G_{47},t)} & \boxed{-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{array}\right] T_{13}$$

$$\frac{dT_{45}}{dt} = (b_{45})^{(9)}T_{44} - \left[\begin{array}{ccc} (b_{45}^{'})^{(9)}\boxed{-(b_{45}^{''})^{(9)}(G_{47},t)} & \boxed{-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{array}\right] T_{14}$$

$$\frac{dT_{46}}{dt} = (b_{46})^{(9)}T_{45} - \left[\begin{array}{ccc} (b_{46}^{'})^{(9)}\boxed{-(b_{46}^{''})^{(9)}(G_{47},t)} & \boxed{-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)} & \boxed{-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)} & \boxed{-(b_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)} & \boxed{-(b_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)} \end{array}\right] T_{15}$$

Where $\boxed{-(b_{44}^{''})^{(9)}(G_{47},t)}$, $\boxed{-(b_{45}^{''})^{(9)}(G_{47},t)}$, $\boxed{-(b_{46}^{''})^{(9)}(G_{47},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)}$, $\boxed{-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)}$ are second detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)}$, $\boxed{-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)}$ are third detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)}$, $\boxed{-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)}$ are fourth detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)}$, $\boxed{-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)}$ are fifth detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)}$, $\boxed{-(b_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)}$ are sixth detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)}$, $\boxed{-(b_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)}$ are seventh detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)}$ are eighth detrition coefficients for category 1, 2 and 3

$\boxed{-(b_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$ are ninth detrition coefficients for category 1, 2 and 3

Where we suppose

$$(a_i)^{(1)}, (a_i^{'})^{(1)}, (a_i^{''})^{(1)}, (b_i)^{(1)}, (b_i^{'})^{(1)}, (b_i^{''})^{(1)} > 0, \quad i,j = 13,14,15$$

The functions $(a_i^{''})^{(1)}, (b_i^{''})^{(1)}$ are positive continuous increasing and bounded.

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| **Definition of** $(p_i)^{(1)}$, $(r_i)^{(1)}$:<br>$(a_i'')^{(1)}(T_{14},t) \le (p_i)^{(1)} \le (\hat{A}_{13})^{(1)}$<br>$(b_i'')^{(1)}(G,t) \le (r_i)^{(1)} \le (b_i')^{(1)} \le (\hat{B}_{13})^{(1)}$ | |
| $lim_{T_2\to\infty}(a_i'')^{(1)}(T_{14},t) = (p_i)^{(1)}$<br>$\lim_{G\to\infty}(b_i'')^{(1)}(G,t) = (r_i)^{(1)}$<br>**Definition of** $(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}$ :<br>Where $\boxed{(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}}$ are positive constants and $\boxed{i = 13,14,15}$ | 98 |
| They satisfy Lipschitz condition:<br>$\lvert(a_i'')^{(1)}(T_{14}',t) - (a_i'')^{(1)}(T_{14},t)\rvert \le (\hat{k}_{13})^{(1)}\lvert T_{14} - T_{14}'\rvert e^{-(\hat{M}_{13})^{(1)}t}$<br>$\lvert(b_i'')^{(1)}(G',t) - (b_i'')^{(1)}(G,t)\rvert < (\hat{k}_{13})^{(1)}\lVert G - G'\rVert e^{-(\hat{M}_{13})^{(1)}t}$ | 99 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(1)}(T_{14}',t)$ and $(a_i'')^{(1)}(T_{14},t)$ . $(T_{14}',t)$ and $(T_{14},t)$ are points belonging to the interval $\left[(\hat{k}_{13})^{(1)}, (\hat{M}_{13})^{(1)}\right]$. It is to be noted that $(a_i'')^{(1)}(T_{14},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{13})^{(1)} = 1$ then the function $(a_i'')^{(1)}(T_{14},t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}$ :<br>$(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}$, are positive constants<br>$\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}} < 1$ | 100 |
| **Definition of** $(\hat{P}_{13})^{(1)}, (\hat{Q}_{13})^{(1)}$ :<br>There exists two constants $(\hat{P}_{13})^{(1)}$ and $(\hat{Q}_{13})^{(1)}$ which together With $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}, (\hat{A}_{13})^{(1)}$ and $(\hat{B}_{13})^{(1)}$ and the constants $(a_i)^{(1)}, (a_i')^{(1)}, (b_i)^{(1)}, (b_i')^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}, i = 13,14,15,$ satisfy the inequalities<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(a_i)^{(1)} + (a_i')^{(1)} + (\hat{A}_{13})^{(1)} + (\hat{P}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(b_i)^{(1)} + (b_i')^{(1)} + (\hat{B}_{13})^{(1)} + (\hat{Q}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$ | 101 |
| Where we suppose | |
| $(a_i)^{(2)}, (a_i')^{(2)}, (a_i'')^{(2)}, (b_i)^{(2)}, (b_i')^{(2)}, (b_i'')^{(2)} > 0, \quad i,j = 16,17,18$ | |

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|---|---|
| The functions $(a_i'')^{(2)}, (b_i'')^{(2)}$ are positive continuous increasing and bounded. | |
| **Definition of** $(p_i)^{(2)}$, $(r_i)^{(2)}$: | |
| $(a_i'')^{(2)}(T_{17},t) \le (p_i)^{(2)} \le (\hat{A}_{16})^{(2)}$ | 102 |
| $(b_i'')^{(2)}(G_{19},t) \le (r_i)^{(2)} \le (b_i')^{(2)} \le (\hat{B}_{16})^{(2)}$ | 103 |
| $\lim_{T_2\to\infty} (a_i'')^{(2)}(T_{17},t) = (p_i)^{(2)}$ | 104 |
| $\lim_{G\to\infty} (b_i'')^{(2)}((G_{19}),t) = (r_i)^{(2)}$ | 105 |
| **Definition of** $(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}$ :<br>Where $(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}, (p_i)^{(2)}, (r_i)^{(2)}$ are positive constants and $i = 16,17,18$ | 106 |
| They satisfy Lipschitz condition: | |
| $\lvert(a_i'')^{(2)}(T_{17}',t) - (a_i'')^{(2)}(T_{17},t)\rvert \le (\hat{k}_{16})^{(2)}\lvert T_{17} - T_{17}'\rvert e^{-(\hat{M}_{16})^{(2)}t}$ | 107 |
| $\lvert(b_i'')^{(2)}((G_{19})',t) - (b_i'')^{(2)}((G_{19}),t)\rvert < (\hat{k}_{16})^{(2)}\lvert\lvert(G_{19}) - (G_{19})'\rvert\rvert e^{-(\hat{M}_{16})^{(2)}t}$ | 108 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(2)}(T_{17}',t)$ and $(a_i'')^{(2)}(T_{17},t)$ . $(T_{17}',t)$ and $(T_{17},t)$ are points belonging to the interval $[(\hat{k}_{16})^{(2)}, (\hat{M}_{16})^{(2)}]$ . It is to be noted that $(a_i'')^{(2)}(T_{17},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{16})^{(2)} = 1$ then the function $(a_i'')^{(2)}(T_{17},t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$ : | |
| $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$, are positive constants<br>$\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}, \frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ | 109 |
| **Definition of** $(\hat{P}_{13})^{(2)}, (\hat{Q}_{13})^{(2)}$ :<br>There exists two constants $(\hat{P}_{16})^{(2)}$ and $(\hat{Q}_{16})^{(2)}$ which together with $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}, (\hat{A}_{16})^{(2)} and (\hat{B}_{16})^{(2)}$ and the constants $(a_i)^{(2)}, (a_i')^{(2)}, (b_i)^{(2)}, (b_i')^{(2)}, (p_i)^{(2)}, (r_i)^{(2)}, i = 16,17,18,$<br>satisfy the inequalities | |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(a_i)^{(2)} + (a_i')^{(2)} + (\hat{A}_{16})^{(2)} + (\hat{P}_{16})^{(2)}(\hat{k}_{16})^{(2)}] < 1$ | 110 |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(b_i)^{(2)} + (b_i')^{(2)} + (\hat{B}_{16})^{(2)} + (\hat{Q}_{16})^{(2)}(\hat{k}_{16})^{(2)}] < 1$ | 111 |

| | |
|---|---|
| Where we suppose | |
| $(a_i)^{(3)}, (a_i^{'})^{(3)}, (a_i^{''})^{(3)}, (b_i)^{(3)}, (b_i^{'})^{(3)}, (b_i^{''})^{(3)} > 0, \quad i,j = 20,21,22$<br><br>The functions $(a_i^{''})^{(3)}, (b_i^{''})^{(3)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(3)}$, $(\mathrm{r_i})^{(3)}$:<br><br>$(a_i^{''})^{(3)}(T_{21}, t) \leq (p_i)^{(3)} \leq (\hat{A}_{20})^{(3)}$<br><br>$(b_i^{''})^{(3)}(G_{23}, t) \leq (r_i)^{(3)} \leq (b_i^{'})^{(3)} \leq (\hat{B}_{20})^{(3)}$ | 112 |
| $lim_{T_2 \to \infty} (a_i^{''})^{(3)} (T_{21}, t) = (p_i)^{(3)}$<br><br>$\lim_{G \to \infty} (b_i^{''})^{(3)} (G_{23}, t) = (r_i)^{(3)}$<br><br>**Definition of** $(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}$ :<br><br>Where $\boxed{(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}}$ are positive constants and $\boxed{i = 20,21,22}$ | 113 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(3)}(T_{21}^{'}, t) - (a_i^{''})^{(3)}(T_{21}, t)\rvert \leq (\hat{k}_{20})^{(3)} \lvert T_{21} - T_{21}^{'}\rvert e^{-(\hat{M}_{20})^{(3)} t}$<br><br>$\lvert(b_i^{''})^{(3)}(G_{23}^{'}, t) - (b_i^{''})^{(3)}(G_{23}, t)\rvert < (\hat{k}_{20})^{(3)} \lvert\lvert G_{23} - G_{23}^{'}\rvert\rvert e^{-(\hat{M}_{20})^{(3)} t}$ | 114 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(3)}(T_{21}^{'}, t)$ and$(a_i^{''})^{(3)}(T_{21}, t)$ . $(T_{21}^{'}, t)$ And $(T_{21}, t)$ are points belonging to the interval $\left[(\hat{k}_{20})^{(3)}, (\hat{M}_{20})^{(3)}\right]$. It is to be noted that $(a_i^{''})^{(3)}(T_{21}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{20})^{(3)} = 1$ then the function $(a_i^{''})^{(3)}(T_{21}, t)$, the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$ :<br><br>$(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$, are positive constants<br>$\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}, \frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}} < 1$ | 115 |
| There exists two constants There exists two constants $(\hat{P}_{20})^{(3)}$ and $(\hat{Q}_{20})^{(3)}$ which together with $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}, (\hat{A}_{20})^{(3)} and\ (\hat{B}_{20})^{(3)}$ and the constants $(a_i)^{(3)}, (a_i^{'})^{(3)}, (b_i)^{(3)}, (b_i^{'})^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}, i = 20,21,22,$<br>satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(a_i)^{(3)} + (a_i^{'})^{(3)} + (\hat{A}_{20})^{(3)} + (\hat{P}_{20})^{(3)} (\hat{k}_{20})^{(3)}] < 1$<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(b_i)^{(3)} + (b_i^{'})^{(3)} + (\hat{B}_{20})^{(3)} + (\hat{Q}_{20})^{(3)} (\hat{k}_{20})^{(3)}] < 1$ | 116 |

Where we suppose

$(a_i)^{(4)}, (a_i^{'})^{(4)}, (a_i^{''})^{(4)}, (b_i)^{(4)}, (b_i^{'})^{(4)}, (b_i^{''})^{(4)} > 0, \quad i,j = 24,25,26$ (117)

The functions $(a_i^{''})^{(4)}, (b_i^{''})^{(4)}$ are positive continuous increasing and bounded.

**Definition of** $(p_i)^{(4)}$, $(r_i)^{(4)}$:

$$(a_i^{''})^{(4)}(T_{25}, t) \leq (p_i)^{(4)} \leq (\hat{A}_{24})^{(4)}$$

$$(b_i^{''})^{(4)}\big((G_{27}), t\big) \leq (r_i)^{(4)} \leq (b_i^{'})^{(4)} \leq (\hat{B}_{24})^{(4)}$$

$lim_{T_2 \to \infty} (a_i^{''})^{(4)}(T_{25}, t) = (p_i)^{(4)}$ (118)

$\lim_{G \to \infty} (b_i^{''})^{(4)}\big((G_{27}), t\big) = (r_i)^{(4)}$

**Definition of** $(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}$ :

Where $\boxed{(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}, (p_i)^{(4)}, (r_i)^{(4)}}$ are positive constants and $\boxed{i = 24,25,26}$

They satisfy Lipschitz condition: (119)

$$|(a_i^{''})^{(4)}(T_{25}^{'}, t) - (a_i^{''})^{(4)}(T_{25}, t)| \leq (\hat{k}_{24})^{(4)} |T_{25} - T_{25}^{'}| e^{-(\hat{M}_{24})^{(4)} t}$$

$$|(b_i^{''})^{(4)}((G_{27})^{'}, t) - (b_i^{''})^{(4)}\big((G_{27}), t\big)| < (\hat{k}_{24})^{(4)} ||(G_{27}) - (G_{27})^{'}|| e^{-(\hat{M}_{24})^{(4)} t}$$

With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(4)}(T_{25}^{'}, t)$ and $(a_i^{''})^{(4)}(T_{25}, t)$ . $(T_{25}^{'}, t)$ and $(T_{25}, t)$ are points belonging to the interval $\big[(\hat{k}_{24})^{(4)}, (\hat{M}_{24})^{(4)}\big]$. It is to be noted that $(a_i^{''})^{(4)}(T_{25}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{24})^{(4)} = 1$ then the function $(a_i^{''})^{(4)}(T_{25}, t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

**Definition of** $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$ : (120)

$(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$, are positive constants

$$\frac{(a_i)^{(4)}}{(\hat{M}_{24})^{(4)}}, \frac{(b_i)^{(4)}}{(\hat{M}_{24})^{(4)}} < 1$$

**Definition of** $(\hat{P}_{24})^{(4)}, (\hat{Q}_{24})^{(4)}$ : (121)

There exists two constants $(\hat{P}_{24})^{(4)}$ and $(\hat{Q}_{24})^{(4)}$ which together with $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}, (\hat{A}_{24})^{(4)} and\ (\hat{B}_{24})^{(4)}$ and the constants $(a_i)^{(4)}, (a_i^{'})^{(4)}, (b_i)^{(4)}, (b_i^{'})^{(4)}, (p_i)^{(4)}, (r_i)^{(4)}, i = 24,25,26$, satisfy the inequalities

$$\frac{1}{(\hat{M}_{24})^{(4)}}[(a_i)^{(4)} + (a_i^{'})^{(4)} + (\hat{A}_{24})^{(4)} + (\hat{P}_{24})^{(4)} (\hat{k}_{24})^{(4)}] < 1$$

| | |
|---|---|
| $\frac{1}{(\hat{M}_{24})^{(4)}}[\ (b_i)^{(4)} + (b_i^{'})^{(4)} +\ (\hat{B}_{24})^{(4)} +\ (\hat{Q}_{24})^{(4)}\ (\hat{k}_{24})^{(4)}] < 1$ | |
| Where we suppose | |
| $(a_i)^{(5)}, (a_i^{'})^{(5)}, (a_i^{''})^{(5)}, (b_i)^{(5)}, (b_i^{'})^{(5)}, (b_i^{''})^{(5)} > 0, \quad i,j = 28,29,30$<br><br>The functions $(a_i^{''})^{(5)}, (b_i^{''})^{(5)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(5)}$, $(r_i)^{(5)}$:<br><br>$(a_i^{''})^{(5)}(T_{29}, t) \leq (p_i)^{(5)} \leq (\hat{A}_{28})^{(5)}$<br><br>$(b_i^{''})^{(5)}\big((G_{31}), t\big) \leq\ (r_i)^{(5)} \leq (b_i^{'})^{(5)} \leq (\hat{B}_{28})^{(5)}$ | 122 |
| $lim_{T_2 \to \infty}(a_i^{''})^{(5)}\,(T_{29}, t) = (p_i)^{(5)}$<br><br>$\lim_{G \to \infty}(b_i^{''})^{(5)}\,(G_{31}, t) =\ (r_i)^{(5)}$<br><br>**Definition of** $(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}$ :<br><br>Where $\boxed{(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}, (p_i)^{(5)},\ (r_i)^{(5)}}$ are positive constants and $\boxed{i = 28,29,30}$ | 123 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(5)}(T_{29}^{'}, t) - (a_i^{''})^{(5)}(T_{29}, t)\rvert \leq (\hat{k}_{28})^{(5)}\lvert T_{29} - T_{29}^{'}\rvert e^{-(\hat{M}_{28})^{(5)}t}$<br><br>$\lvert(b_i^{''})^{(5)}((G_{31})^{'}, t) - (b_i^{''})^{(5)}\big((G_{31}), t\big)\rvert < (\hat{k}_{28})^{(5)}\lvert\lvert(G_{31}) - (G_{31})^{'}\rvert\rvert e^{-(\hat{M}_{28})^{(5)}t}$ | 124 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(5)}(T_{29}^{'}, t)$ and $(a_i^{''})^{(5)}(T_{29}, t)$ . $(T_{29}^{'}, t)$ and $(T_{29}, t)$ are points belonging to the interval $\big[(\hat{k}_{28})^{(5)}, (\hat{M}_{28})^{(5)}\big]$. It is to be noted that $(a_i^{''})^{(5)}(T_{29}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{28})^{(5)} = 1$ then the function $(a_i^{''})^{(5)}(T_{29}, t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}$ :<br><br>$(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}$, are positive constants<br><br>$\frac{(a_i)^{(5)}}{(\hat{M}_{28})^{(5)}}\ , \frac{(b_i)^{(5)}}{(\hat{M}_{28})^{(5)}} < 1$ | 125 |
| **Definition of** $(\hat{P}_{28})^{(5)}, (\hat{Q}_{28})^{(5)}$ :<br><br>There exists two constants $(\hat{P}_{28})^{(5)}$ and $(\hat{Q}_{28})^{(5)}$ which together with $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}, (\hat{A}_{28})^{(5)} and\ (\hat{B}_{28})^{(5)}$ and the constants $(a_i)^{(5)}, (a_i^{'})^{(5)}, (b_i)^{(5)}, (b_i^{'})^{(5)}, (p_i)^{(5)},\ (r_i)^{(5)}, i = 28,29,30,$ satisfy the inequalities | 126 |

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| $\frac{1}{(\hat{M}_{28})^{(5)}}[(a_i)^{(5)} + (a_i^{'})^{(5)} + (\hat{A}_{28})^{(5)} + (\hat{P}_{28})^{(5)}(\hat{k}_{28})^{(5)}] < 1$<br><br>$\frac{1}{(\hat{M}_{28})^{(5)}}[(b_i)^{(5)} + (b_i^{'})^{(5)} + (\hat{B}_{28})^{(5)} + (\hat{Q}_{28})^{(5)}(\hat{k}_{28})^{(5)}] < 1$ | |
| Where we suppose | |
| $(a_i)^{(6)}, (a_i^{'})^{(6)}, (a_i^{''})^{(6)}, (b_i)^{(6)}, (b_i^{'})^{(6)}, (b_i^{''})^{(6)} > 0, \quad i,j = 32,33,34$<br><br>The functions $(a_i^{''})^{(6)}, (b_i^{''})^{(6)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(6)}$, $(r_i)^{(6)}$:<br><br>$(a_i^{''})^{(6)}(T_{33}, t) \leq (p_i)^{(6)} \leq (\hat{A}_{32})^{(6)}$<br><br>$(b_i^{''})^{(6)}((G_{35}), t) \leq (r_i)^{(6)} \leq (b_i^{'})^{(6)} \leq (\hat{B}_{32})^{(6)}$ | 127 |
| $lim_{T_2 \to \infty}(a_i^{''})^{(6)}(T_{33}, t) = (p_i)^{(6)}$<br><br>$\lim_{G \to \infty}(b_i^{''})^{(6)}((G_{35}), t) = (r_i)^{(6)}$<br><br>**Definition of** $(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}$ :<br><br>Where $\boxed{(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}, (p_i)^{(6)}, (r_i)^{(6)}}$ are positive constants and $\boxed{i = 32,33,34}$ | 128 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(6)}(T_{33}^{'}, t) - (a_i^{''})^{(6)}(T_{33}, t)\rvert \leq (\hat{k}_{32})^{(6)}\lvert T_{33} - T_{33}^{'}\rvert e^{-(\hat{M}_{32})^{(6)}t}$<br><br>$\lvert(b_i^{''})^{(6)}((G_{35})^{'}, t) - (b_i^{''})^{(6)}((G_{35}), t)\rvert < (\hat{k}_{32})^{(6)}\lvert\lvert(G_{35}) - (G_{35})^{'}\rvert\rvert e^{-(\hat{M}_{32})^{(6)}t}$ | |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(6)}(T_{33}^{'}, t)$ and$(a_i^{''})^{(6)}(T_{33}, t)$ . $(T_{33}^{'}, t)$ and $(T_{33}, t)$ are points belonging to the interval $[(\hat{k}_{32})^{(6)}, (\hat{M}_{32})^{(6)}]$. It is to be noted that $(a_i^{''})^{(6)}(T_{33}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{32})^{(6)} = 1$ then the function $(a_i^{''})^{(6)}(T_{33}, t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}$ :<br><br>$(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}$, are positive constants<br><br>$\frac{(a_i)^{(6)}}{(\hat{M}_{32})^{(6)}}, \frac{(b_i)^{(6)}}{(\hat{M}_{32})^{(6)}} < 1$ | 129 |
| **Definition of** $(\hat{P}_{32})^{(6)}, (\hat{Q}_{32})^{(6)}$ :<br><br>There exists two constants $(\hat{P}_{32})^{(6)}$ and $(\hat{Q}_{32})^{(6)}$ which together with $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}, (\hat{A}_{32})^{(6)} and\ (\hat{B}_{32})^{(6)}$ and the constants | 130 |

$(a_i)^{(6)}, (a_i^{'})^{(6)}, (b_i)^{(6)}, (b_i^{'})^{(6)}, (p_i)^{(6)}, \ (r_i)^{(6)}, i = 32,33,34,$
satisfy the inequalities

$$\frac{1}{(\hat{M}_{32})^{(6)}}[(a_i)^{(6)} + (a_i^{'})^{(6)} + (\hat{A}_{32})^{(6)} + (\hat{P}_{32})^{(6)} (\hat{k}_{32})^{(6)}] < 1$$

$$\frac{1}{(\hat{M}_{32})^{(6)}}[(b_i)^{(6)} + (b_i^{'})^{(6)} + (\hat{B}_{32})^{(6)} + (\hat{Q}_{32})^{(6)} (\hat{k}_{32})^{(6)}] < 1$$

Where we suppose

(KK) $(a_i)^{(7)}, (a_i^{'})^{(7)}, (a_i^{''})^{(7)}, (b_i)^{(7)}, (b_i^{'})^{(7)}, (b_i^{''})^{(7)} > 0, \quad i, j = 36,37,38$ 131

(LL) The functions $(a_i^{''})^{(7)}, (b_i^{''})^{(7)}$ are positive continuous increasing and bounded.

**Definition of** $(p_i)^{(7)}, \ (r_i)^{(7)}$:

$$(a_i^{''})^{(7)}(T_{37}, t) \leq (p_i)^{(7)} \leq (\hat{A}_{36})^{(7)}$$

$$(b_i^{''})^{(7)}(G_{39}, t) \leq (r_i)^{(7)} \leq (b_i^{'})^{(7)} \leq (\hat{B}_{36})^{(7)}$$

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*(MM)* $\lim_{T_2 \to \infty} (a_i^{''})^{(7)}(T_{37}, t) = (p_i)^{(7)}$

*(NN)*

$\lim_{G \to \infty} (b_i^{''})^{(7)}((G_{39}), t) = (r_i)^{(7)}$

**Definition of** $(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}$ :

Where $\boxed{(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}, (p_i)^{(7)}, \ (r_i)^{(7)}}$ are positive constants and $\boxed{i = 36,37,38}$

They satisfy Lipschitz condition: 133

$$|(a_i^{''})^{(7)}(T_{37}^{'}, t) - (a_i^{''})^{(7)}(T_{37}, t)| \leq (\hat{k}_{36})^{(7)} |T_{37} - T_{37}^{'}| e^{-(\hat{M}_{36})^{(7)} t}$$

$$|(b_i^{''})^{(7)}((G_{39})^{'}, t) - (b_i^{''})^{(7)}((G_{39}), t)| < (\hat{k}_{36})^{(7)} ||(G_{39}) - (G_{39})^{'}|| e^{-(\hat{M}_{36})^{(7)} t}$$

With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(7)}(T_{37}^{'}, t)$ and $(a_i^{''})^{(7)}(T_{37}, t)$ . $(T_{37}^{'}, t)$ and $(T_{37}, t)$ are points belonging to the interval $[(\hat{k}_{36})^{(7)}, (\hat{M}_{36})^{(7)}]$ . It is to be noted that $(a_i^{''})^{(7)}(T_{37}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{36})^{(7)} = 1$ then the function $(a_i^{''})^{(7)}(T_{37}, t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

**Definition of** $(\hat{M}_{36})^{(7)}, (\hat{k}_{36})^{(7)}$ : 134

(OO) $(\hat{M}_{36})^{(7)}, (\hat{k}_{36})^{(7)},$ are positive constants

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| $\frac{(a_i)^{(7)}}{(\hat{M}_{36})^{(7)}}, \frac{(b_i)^{(7)}}{(\hat{M}_{36})^{(7)}} < 1$ | |
| **Definition of** $(\hat{P}_{36})^{(7)}, (\hat{Q}_{36})^{(7)}$ :<br><br>(PP) There exists two constants $(\hat{P}_{36})^{(7)}$ and $(\hat{Q}_{36})^{(7)}$ which together with $(\hat{M}_{36})^{(7)}, (\hat{k}_{36})^{(7)}, (\hat{A}_{36})^{(7)} and\ (\hat{B}_{36})^{(7)}$ and the constants $(a_i)^{(7)}, (a_i^{'})^{(7)}, (b_i)^{(7)}, (b_i^{'})^{(7)}, (p_i)^{(7)}, (r_i)^{(7)}, i = 36,37,38$, satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{36})^{(7)}}[(a_i)^{(7)} + (a_i^{'})^{(7)} + (\hat{A}_{36})^{(7)} + (\hat{P}_{36})^{(7)} (\hat{k}_{36})^{(7)}] < 1$<br><br>$\frac{1}{(\hat{M}_{36})^{(7)}}[(b_i)^{(7)} + (b_i^{'})^{(7)} + (\hat{B}_{36})^{(7)} + (\hat{Q}_{36})^{(7)} (\hat{k}_{36})^{(7)}] < 1$ | 135 |
| Where we suppose | |
| $(a_i)^{(8)}, (a_i^{'})^{(8)}, (a_i^{''})^{(8)}, (b_i)^{(8)}, (b_i^{'})^{(8)}, (b_i^{''})^{(8)} > 0, \quad i,j = 40,41,42$ | 136 |
| The functions $(a_i^{''})^{(8)}, (b_i^{''})^{(8)}$ are positive continuous increasing and bounded | |
| **Definition of** $(p_i)^{(8)}, (r_i)^{(8)}$: | 137 |
| $(a_i^{''})^{(8)}(T_{41}, t) \leq (p_i)^{(8)} \leq (\hat{A}_{40})^{(8)}$ | 138 |
| $(b_i^{''})^{(8)}((G_{43}), t) \leq (r_i)^{(8)} \leq (b_i^{'})^{(8)} \leq (\hat{B}_{40})^{(8)}$ | 139 |
| $\lim_{T_2 \to \infty} (a_i^{''})^{(8)}(T_{41}, t) = (p_i)^{(8)}$ | 140 |
| $\lim_{G \to \infty} (b_i^{''})^{(8)}((G_{43}), t) = (r_i)^{(8)}$ | 141 |
| **Definition of** $(\hat{A}_{40})^{(8)}, (\hat{B}_{40})^{(8)}$ :<br><br>Where $\boxed{(\hat{A}_{40})^{(8)}, (\hat{B}_{40})^{(8)}, (p_i)^{(8)}, (r_i)^{(8)}}$ are positive constants and $\boxed{i = 40,41,42}$ | |
| They satisfy Lipschitz condition: | |
| $\lvert(a_i^{''})^{(8)}(T_{41}^{'}, t) - (a_i^{''})^{(8)}(T_{41}, t)\rvert \leq (\hat{k}_{40})^{(8)} \lvert T_{41} - T_{41}^{'}\rvert e^{-(\hat{M}_{40})^{(8)} t}$ | 142 |
| $\lvert(b_i^{''})^{(8)}((G_{43})', t) - (b_i^{''})^{(8)}((G_{43}), t)\rvert < (\hat{k}_{40})^{(8)} \lvert\lvert(G_{43}) - (G_{43})'\rvert\rvert e^{-(\hat{M}_{40})^{(8)} t}$ | 143 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(8)}(T_{41}^{'}, t)$ and $(a_i^{''})^{(8)}(T_{41}, t)$ . $(T_{41}^{'}, t)$ and $(T_{41}, t)$ are points belonging to the interval $[(\hat{k}_{40})^{(8)}, (\hat{M}_{40})^{(8)}]$. It is to be noted that $(a_i^{''})^{(8)}(T_{41}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{40})^{(8)} = 1$ then the function $(a_i^{''})^{(8)}(T_{41}, t)$ , the **first augmentation coefficient** attributable to the system, would | |

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| be absolutely continuous. | |
| **Definition of** $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}$ : | |
| $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)},$ are positive constants | |
| $\frac{(a_i)^{(8)}}{(\hat{M}_{40})^{(8)}}$ , $\frac{(b_i)^{(8)}}{(\hat{M}_{40})^{(8)}} < 1$ | 144 |
| **Definition of** $(\hat{P}_{40})^{(8)}, (\hat{Q}_{40})^{(8)}$ :<br>There exists two constants $(\hat{P}_{40})^{(8)}$ and $(\hat{Q}_{40})^{(8)}$ which together with $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}, (\hat{A}_{40})^{(8)}$ $(\hat{B}_{40})^{(8)}$ and the constants $(a_i)^{(8)}, (a_i')^{(8)}, (b_i)^{(8)}, (b_i')^{(8)}, (p_i)^{(8)}, (r_i)^{(8)}, i = 40,41,42,$<br>Satisfy the inequalities | |
| $\frac{1}{(\hat{M}_{40})^{(8)}}[(a_i)^{(8)} + (a_i')^{(8)} + (\hat{A}_{40})^{(8)} + (\hat{P}_{40})^{(8)} (\hat{k}_{40})^{(8)}] < 1$ | 145 |
| $\frac{1}{(\hat{M}_{40})^{(8)}}[(b_i)^{(8)} + (b_i')^{(8)} + (\hat{B}_{40})^{(8)} + (\hat{Q}_{40})^{(8)} (\hat{k}_{40})^{(8)}] < 1$ | 146 |
| Where we suppose | |
| $(a_i)^{(9)}, (a_i')^{(9)}, (a_i'')^{(9)}, (b_i)^{(9)}, (b_i')^{(9)}, (b_i'')^{(9)} > 0, \quad i,j = 44,45,46$<br>The functions $(a_i'')^{(9)}, (b_i'')^{(9)}$ are positive continuous increasing and bounded.<br>**Definition of** $(p_i)^{(9)}$, $(r_i)^{(9)}$:<br>$(a_i'')^{(9)}(T_{45}, t) \le (p_i)^{(9)} \le (\hat{A}_{44})^{(9)}$<br>$(b_i'')^{(9)}(G_{47}, t) \le (r_i)^{(9)} \le (b_i')^{(9)} \le (\hat{B}_{44})^{(9)}$ | 146 A |
| $\lim_{T_2 \to \infty} (a_i'')^{(9)}(T_{45}, t) = (p_i)^{(9)}$<br>$\lim_{G \to \infty} (b_i'')^{(9)}(G_{47}, t) = (r_i)^{(9)}$<br>**Definition of** $(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}$ :<br>Where $\boxed{(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}, (p_i)^{(9)}, (r_i)^{(9)}}$ are positive constants and $\boxed{i = 44,45,46}$ | |
| They satisfy Lipschitz condition:<br>$\lvert(a_i'')^{(9)}(T_{45}', t) - (a_i'')^{(9)}(T_{45}, t)\rvert \le (\hat{k}_{44})^{(9)} \lvert T_{45} - T_{45}' \rvert e^{-(\hat{M}_{44})^{(9)} t}$<br>$\lvert(b_i'')^{(9)}((G_{47})', t) - (b_i'')^{(9)}((G_{47}), t)\rvert < (\hat{k}_{44})^{(9)} \lvert\lvert(G_{47}) - (G_{47})'\rvert\rvert e^{-(\hat{M}_{44})^{(9)} t}$ | |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(9)}(T_{45}', t)$ | |

and $(a_i'')^{(9)}(T_{45},t)$ . $(T_{45}',t)$ and $(T_{45},t)$ are points belonging to the interval $\left[(\hat{k}_{44})^{(9)},(\widehat{M}_{44})^{(9)}\right]$. It is to be noted that $(a_i'')^{(9)}(T_{45},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\widehat{M}_{44})^{(9)}=1$ then the function $(a_i'')^{(9)}(T_{45},t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

**Definition of** $(\widehat{M}_{44})^{(9)},(\hat{k}_{44})^{(9)}$ :

$(\widehat{M}_{44})^{(9)},(\hat{k}_{44})^{(9)},$ are positive constants

$$\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\ ,\frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}<1$$

**Definition of** $(\hat{P}_{44})^{(9)},(\hat{Q}_{44})^{(9)}$ :

There exists two constants $(\hat{P}_{44})^{(9)}$ and $(\hat{Q}_{44})^{(9)}$ which together with $(\widehat{M}_{44})^{(9)},(\hat{k}_{44})^{(9)},(\hat{A}_{44})^{(9)} and\ (\hat{B}_{44})^{(9)}$ and the constants $(a_i)^{(9)},(a_i')^{(9)},(b_i)^{(9)},(b_i')^{(9)},(p_i)^{(9)},\ (r_i)^{(9)},i=44,45,46,$ satisfy the inequalities

$$\frac{1}{(\widehat{M}_{44})^{(9)}}[(a_i)^{(9)}+(a_i')^{(9)}+\ (\hat{A}_{44})^{(9)}+\ (\hat{P}_{44})^{(9)}\ (\hat{k}_{44})^{(9)}]<1$$

$$\frac{1}{(\widehat{M}_{44})^{(9)}}[\ (b_i)^{(9)}+(b_i')^{(9)}+\ (\hat{B}_{44})^{(9)}+\ (\hat{Q}_{44})^{(9)}\ \ (\hat{k}_{44})^{(9)}]<1$$

**Theorem 1:** if the conditions above are fulfilled, there exists a solution satisfying the conditions 147

Definition of $G_i(0)\ ,T_i(0)$ :

$$G_i(t)\le\ \left(\hat{P}_{13}\right)^{(1)}e^{(\widehat{M}_{13})^{(1)}t}\ ,\ \boxed{G_i(0)=G_i^0>0}$$

$$T_i(t)\le\ (\hat{Q}_{13})^{(1)}e^{(\widehat{M}_{13})^{(1)}t}\ ,\ \boxed{T_i(0)=T_i^0>0}$$

**Theorem 2 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions 148

**Definition of** $G_i(0)\ ,T_i(0)$

$$G_i(t)\le\ (\hat{P}_{16})^{(2)}e^{(\widehat{M}_{16})^{(2)}t}\ ,\quad G_i(0)=G_i^0>0$$

$$T_i(t)\le\ (\hat{Q}_{16})^{(2)}e^{(\widehat{M}_{16})^{(2)}t}\ ,\quad T_i(0)=T_i^0>0$$

**Theorem 3 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions 149

$$G_i(t)\le\ (\hat{P}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}\ ,\quad G_i(0)=G_i^0>0$$

$$T_i(t)\le\ (\hat{Q}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}\ ,\quad T_i(0)=T_i^0>0$$

**Theorem 4 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions 150

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| **Definition of** $G_i(0)\,, T_i(0)$ :<br>$G_i(t) \le \left(\hat{P}_{24}\right)^{(4)} e^{(\hat{M}_{24})^{(4)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{24}\right)^{(4)} e^{(\hat{M}_{24})^{(4)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | |
| **Theorem 5 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0)\,, T_i(0)$ :<br>$G_i(t) \le \left(\hat{P}_{28}\right)^{(5)} e^{(\hat{M}_{28})^{(5)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{28}\right)^{(5)} e^{(\hat{M}_{28})^{(5)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 151 |
| **Theorem 6 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0)\,, T_i(0)$ :<br>$G_i(t) \le \left(\hat{P}_{32}\right)^{(6)} e^{(\hat{M}_{32})^{(6)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{32}\right)^{(6)} e^{(\hat{M}_{32})^{(6)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 152 |
| **Theorem 7:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0)\,, T_i(0)$ :<br>$G_i(t) \le \left(\hat{P}_{36}\right)^{(7)} e^{(\hat{M}_{36})^{(7)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{36}\right)^{(7)} e^{(\hat{M}_{36})^{(7)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 |
| **Theorem 8:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0)\,, T_i(0)$ :<br>$G_i(t) \le \left(\hat{P}_{40}\right)^{(8)} e^{(\hat{M}_{40})^{(8)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \le \left(\hat{Q}_{40}\right)^{(8)} e^{(\hat{M}_{40})^{(8)}t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 A |
| **Theorem 9:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0)\,, T_i(0)$ :<br>$G_i(t) \le \left(\hat{P}_{44}\right)^{(9)} e^{(\hat{M}_{44})^{(9)}t}$ , $\boxed{G_i(0) = G_i^0 > 0}$ | 153 B |

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| $T_i(t) \leq (\hat{Q}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | |
| **Proof:** Consider operator $\mathcal{A}^{(1)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 154 |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{13})^{(1)}$ , $T_i^0 \leq (\hat{Q}_{13})^{(1)}$, | 155 |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ | 156 |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ | 157 |
| By<br>$\bar{G}_{13}(t) = G_{13}^0 + \int_0^t \left[ (a_{13})^{(1)} G_{14}(s_{(13)}) - \left( (a_{13}')^{(1)} + a_{13}'')^{(1)} (T_{14}(s_{(13)}), s_{(13)}) \right) G_{13}(s_{(13)}) \right] ds_{(13)}$ | 158 |
| $\bar{G}_{14}(t) = G_{14}^0 + \int_0^t \left[ (a_{14})^{(1)} G_{13}(s_{(13)}) - \left( (a_{14}')^{(1)} + (a_{14}'')^{(1)} (T_{14}(s_{(13)}), s_{(13)}) \right) G_{14}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{G}_{15}(t) = G_{15}^0 + \int_0^t \left[ (a_{15})^{(1)} G_{14}(s_{(13)}) - \left( (a_{15}')^{(1)} + (a_{15}'')^{(1)} (T_{14}(s_{(13)}), s_{(13)}) \right) G_{15}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{T}_{13}(t) = T_{13}^0 + \int_0^t \left[ (b_{13})^{(1)} T_{14}(s_{(13)}) - \left( (b_{13}')^{(1)} - (b_{13}'')^{(1)} (G(s_{(13)}), s_{(13)}) \right) T_{13}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{T}_{14}(t) = T_{14}^0 + \int_0^t \left[ (b_{14})^{(1)} T_{13}(s_{(13)}) - \left( (b_{14}')^{(1)} - (b_{14}'')^{(1)} (G(s_{(13)}), s_{(13)}) \right) T_{14}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{T}_{15}(t) = T_{15}^0 + \int_0^t \left[ (b_{15})^{(1)} T_{14}(s_{(13)}) - \left( (b_{15}')^{(1)} - (b_{15}'')^{(1)} (G(s_{(13)}), s_{(13)}) \right) T_{15}(s_{(13)}) \right] ds_{(13)}$<br>Where $s_{(13)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(2)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 159 |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{16})^{(2)}$ , $T_i^0 \leq (\hat{Q}_{16})^{(2)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ | |
| By<br>$\bar{G}_{16}(t) = G_{16}^0 + \int_0^t \left[ (a_{16})^{(2)} G_{17}(s_{(16)}) - \left( (a_{16}')^{(2)} + a_{16}'')^{(2)} (T_{17}(s_{(16)}), s_{(16)}) \right) G_{16}(s_{(16)}) \right] ds_{(16)}$ | 160 |
| $\bar{G}_{17}(t) = G_{17}^0 + \int_0^t \left[ (a_{17})^{(2)} G_{16}(s_{(16)}) - \left( (a_{17}')^{(2)} + (a_{17}'')^{(2)} (T_{17}(s_{(16)}), s_{(17)}) \right) G_{17}(s_{(16)}) \right] ds_{(16)}$ | |

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| $\bar{G}_{18}(t) = G_{18}^0 + \int_0^t \left[(a_{18})^{(2)} G_{17}(s_{(16)}) - \left((a_{18}')^{(2)} + (a_{18}'')^{(2)}(T_{17}(s_{(16)}), s_{(16)})\right) G_{18}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{T}_{16}(t) = T_{16}^0 + \int_0^t \left[(b_{16})^{(2)} T_{17}(s_{(16)}) - \left((b_{16}')^{(2)} - (b_{16}'')^{(2)}(G_{19}(s_{(16)}), s_{(16)})\right) T_{16}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{T}_{17}(t) = T_{17}^0 + \int_0^t \left[(b_{17})^{(2)} T_{16}(s_{(16)}) - \left((b_{17}')^{(2)} - (b_{17}'')^{(2)}(G_{19}(s_{(16)}), s_{(16)})\right) T_{17}(s_{(16)})\right] ds_{(16)}$ | |
| $\bar{T}_{18}(t) = T_{18}^0 + \int_0^t \left[(b_{18})^{(2)} T_{17}(s_{(16)}) - \left((b_{18}')^{(2)} - (b_{18}'')^{(2)}(G_{19}(s_{(16)}), s_{(16)})\right) T_{18}(s_{(16)})\right] ds_{(16)}$<br>Where $s_{(16)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(3)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{20})^{(3)}$, $T_i^0 \le (\hat{Q}_{20})^{(3)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ | |
| By<br>$\bar{G}_{20}(t) = G_{20}^0 + \int_0^t \left[(a_{20})^{(3)} G_{21}(s_{(20)}) - \left((a_{20}')^{(3)} + a_{20}'')^{(3)}(T_{21}(s_{(20)}), s_{(20)})\right) G_{20}(s_{(20)})\right] ds_{(20)}$ | 161 |
| $\bar{G}_{21}(t) = G_{21}^0 + \int_0^t \left[(a_{21})^{(3)} G_{20}(s_{(20)}) - \left((a_{21}')^{(3)} + (a_{21}'')^{(3)}(T_{21}(s_{(20)}), s_{(20)})\right) G_{21}(s_{(20)})\right] ds_{(20)}$ | |
| $\bar{G}_{22}(t) = G_{22}^0 + \int_0^t \left[(a_{22})^{(3)} G_{21}(s_{(20)}) - \left((a_{22}')^{(3)} + (a_{22}'')^{(3)}(T_{21}(s_{(20)}), s_{(20)})\right) G_{22}(s_{(20)})\right] ds_{(20)}$ | |
| $\bar{T}_{20}(t) = T_{20}^0 + \int_0^t \left[(b_{20})^{(3)} T_{21}(s_{(20)}) - \left((b_{20}')^{(3)} - (b_{20}'')^{(3)}(G_{23}(s_{(20)}), s_{(20)})\right) T_{20}(s_{(20)})\right] ds_{(20)}$ | |
| $\bar{T}_{21}(t) = T_{21}^0 + \int_0^t \left[(b_{21})^{(3)} T_{20}(s_{(20)}) - \left((b_{21}')^{(3)} - (b_{21}'')^{(3)}(G_{23}(s_{(20)}), s_{(20)})\right) T_{21}(s_{(20)})\right] ds_{(20)}$ | |
| $\bar{T}_{22}(t) = T_{22}^0 + \int_0^t \left[(b_{22})^{(3)} T_{21}(s_{(20)}) - \left((b_{22}')^{(3)} - (b_{22}'')^{(3)}(G_{23}(s_{(20)}), s_{(20)})\right) T_{22}(s_{(20)})\right] ds_{(20)}$<br>Where $s_{(20)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:** Consider operator $\mathcal{A}^{(4)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{24})^{(4)}$, $T_i^0 \le (\hat{Q}_{24})^{(4)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ | |

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| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ | |
| By<br>$\bar{G}_{24}(t) = G_{24}^0 + \int_0^t \left[ (a_{24})^{(4)} G_{25}(s_{(24)}) - \left( (a'_{24})^{(4)} + a''_{24})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{24}(s_{(24)}) \right] ds_{(24)}$ | 162 |
| $\bar{G}_{25}(t) = G_{25}^0 + \int_0^t \left[ (a_{25})^{(4)} G_{24}(s_{(24)}) - \left( (a'_{25})^{(4)} + (a''_{25})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{25}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{G}_{26}(t) = G_{26}^0 + \int_0^t \left[ (a_{26})^{(4)} G_{25}(s_{(24)}) - \left( (a'_{26})^{(4)} + (a''_{26})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{26}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{24}(t) = T_{24}^0 + \int_0^t \left[ (b_{24})^{(4)} T_{25}(s_{(24)}) - \left( (b'_{24})^{(4)} - (b''_{24})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{24}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{25}(t) = T_{25}^0 + \int_0^t \left[ (b_{25})^{(4)} T_{24}(s_{(24)}) - \left( (b'_{25})^{(4)} - (b''_{25})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{25}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{26}(t) = T_{26}^0 + \int_0^t \left[ (b_{26})^{(4)} T_{25}(s_{(24)}) - \left( (b'_{26})^{(4)} - (b''_{26})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{26}(s_{(24)}) \right] ds_{(24)}$<br>Where $s_{(24)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:** Consider operator $\mathcal{A}^{(5)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{28})^{(5)}$, $T_i^0 \le (\hat{Q}_{28})^{(5)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| By<br>$\bar{G}_{28}(t) = G_{28}^0 + \int_0^t \left[ (a_{28})^{(5)} G_{29}(s_{(28)}) - \left( (a'_{28})^{(5)} + a''_{28})^{(5)} (T_{29}(s_{(28)}), s_{(28)}) \right) G_{28}(s_{(28)}) \right] ds_{(28)}$ | 163 |
| $\bar{G}_{29}(t) = G_{29}^0 + \int_0^t \left[ (a_{29})^{(5)} G_{28}(s_{(28)}) - \left( (a'_{29})^{(5)} + (a''_{29})^{(5)} (T_{29}(s_{(28)}), s_{(28)}) \right) G_{29}(s_{(28)}) \right] ds_{(28)}$ | |
| $\bar{G}_{30}(t) = G_{30}^0 + \int_0^t \left[ (a_{30})^{(5)} G_{29}(s_{(28)}) - \left( (a'_{30})^{(5)} + (a''_{30})^{(5)} (T_{29}(s_{(28)}), s_{(28)}) \right) G_{30}(s_{(28)}) \right] ds_{(28)}$ | |
| $\bar{T}_{28}(t) = T_{28}^0 + \int_0^t \left[ (b_{28})^{(5)} T_{29}(s_{(28)}) - \left( (b'_{28})^{(5)} - (b''_{28})^{(5)} (G_{31}(s_{(28)}), s_{(28)}) \right) T_{28}(s_{(28)}) \right] ds_{(28)}$ | |
| $\bar{T}_{29}(t) = T_{29}^0 + \int_0^t \left[ (b_{29})^{(5)} T_{28}(s_{(28)}) - \left( (b'_{29})^{(5)} - (b''_{29})^{(5)} (G_{31}(s_{(28)}), s_{(28)}) \right) T_{29}(s_{(28)}) \right] ds_{(28)}$ | |
| $\bar{T}_{30}(t) = T_{30}^0 + \int_0^t \left[ (b_{30})^{(5)} T_{29}(s_{(28)}) - \left( (b'_{30})^{(5)} - (b''_{30})^{(5)} (G_{31}(s_{(28)}), s_{(28)}) \right) T_{30}(s_{(28)}) \right] ds_{(28)}$<br>Where $s_{(28)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:** | |

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| Consider operator $\mathcal{A}^{(6)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{32})^{(6)}$ , $T_i^0 \leq (\hat{Q}_{32})^{(6)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ | |
| By<br>$\bar{G}_{32}(t) = G_{32}^0 + \int_0^t \left[ (a_{32})^{(6)} G_{33}(s_{(32)}) - \left( (a'_{32})^{(6)} + a''_{32})^{(6)} (T_{33}(s_{(32)}), s_{(32)}) \right) G_{32}(s_{(32)}) \right] ds_{(32)}$ | 164 |
| $\bar{G}_{33}(t) = G_{33}^0 + \int_0^t \left[ (a_{33})^{(6)} G_{32}(s_{(32)}) - \left( (a'_{33})^{(6)} + (a''_{33})^{(6)} (T_{33}(s_{(32)}), s_{(32)}) \right) G_{33}(s_{(32)}) \right] ds_{(32)}$ | |
| $\bar{G}_{34}(t) = G_{34}^0 + \int_0^t \left[ (a_{34})^{(6)} G_{33}(s_{(32)}) - \left( (a'_{34})^{(6)} + (a''_{34})^{(6)} (T_{33}(s_{(32)}), s_{(32)}) \right) G_{34}(s_{(32)}) \right] ds_{(32)}$ | |
| $\bar{T}_{32}(t) = T_{32}^0 + \int_0^t \left[ (b_{32})^{(6)} T_{33}(s_{(32)}) - \left( (b'_{32})^{(6)} - (b''_{32})^{(6)} (G_{35}(s_{(32)}), s_{(32)}) \right) T_{32}(s_{(32)}) \right] ds_{(32)}$ | |
| $\bar{T}_{33}(t) = T_{33}^0 + \int_0^t \left[ (b_{33})^{(6)} T_{32}(s_{(32)}) - \left( (b'_{33})^{(6)} - (b''_{33})^{(6)} (G_{35}(s_{(32)}), s_{(32)}) \right) T_{33}(s_{(32)}) \right] ds_{(32)}$ | |
| $\bar{T}_{34}(t) = T_{34}^0 + \int_0^t \left[ (b_{34})^{(6)} T_{33}(s_{(32)}) - \left( (b'_{34})^{(6)} - (b''_{34})^{(6)} (G_{35}(s_{(32)}), s_{(32)}) \right) T_{34}(s_{(32)}) \right] ds_{(32)}$<br>Where $s_{(32)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(7)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{36})^{(7)}$ , $T_i^0 \leq (\hat{Q}_{36})^{(7)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ | |
| By<br>$\bar{G}_{36}(t) = G_{36}^0 + \int_0^t \left[ (a_{36})^{(7)} G_{37}(s_{(36)}) - \left( (a'_{36})^{(7)} + a''_{36})^{(7)} (T_{37}(s_{(36)}), s_{(36)}) \right) G_{36}(s_{(36)}) \right] ds_{(36)}$ | 165 |
| $\bar{G}_{37}(t) = G_{37}^0 + \int_0^t \left[ (a_{37})^{(7)} G_{36}(s_{(36)}) - \left( (a'_{37})^{(7)} + (a''_{37})^{(7)} (T_{37}(s_{(36)}), s_{(36)}) \right) G_{37}(s_{(36)}) \right] ds_{(36)}$ | |
| $\bar{G}_{38}(t) = G_{38}^0 + \int_0^t \left[ (a_{38})^{(7)} G_{37}(s_{(36)}) - \left( (a'_{38})^{(7)} + (a''_{38})^{(7)} (T_{37}(s_{(36)}), s_{(36)}) \right) G_{38}(s_{(36)}) \right] ds_{(36)}$ | |

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| $\bar{T}_{36}(t) = T_{36}^{0} + \int_0^t \left[(b_{36})^{(7)} T_{37}(s_{(36)}) - \left((b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{36}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{37}(t) = T_{37}^{0} + \int_0^t \left[(b_{37})^{(7)} T_{36}(s_{(36)}) - \left((b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{37}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{38}(t) = T_{38}^{0} + \int_0^t \left[(b_{38})^{(7)} T_{37}(s_{(36)}) - \left((b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{38}(s_{(36)})\right] ds_{(36)}$<br>Where $s_{(36)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(8)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \le (\hat{P}_{40})^{(8)}$ , $T_i^0 \le (\hat{Q}_{40})^{(8)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{40})^{(8)} e^{(\hat{M}_{40})^{(8)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{40})^{(8)} e^{(\hat{M}_{40})^{(8)} t}$ | |
| By<br>$\bar{G}_{40}(t) = G_{40}^0 + \int_0^t \left[(a_{40})^{(8)} G_{41}(s_{(40)}) - \left((a_{40}^{'})^{(8)} + a_{40}^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\right) G_{40}(s_{(40)})\right] ds_{(40)}$ | 166 |
| $\bar{G}_{41}(t) = G_{41}^0 + \int_0^t \left[(a_{41})^{(8)} G_{40}(s_{(40)}) - \left((a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\right) G_{41}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{G}_{42}(t) = G_{42}^0 + \int_0^t \left[(a_{42})^{(8)} G_{41}(s_{(40)}) - \left((a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\right) G_{42}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{T}_{40}(t) = T_{40}^{0} + \int_0^t \left[(b_{40})^{(8)} T_{41}(s_{(40)}) - \left((b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}(G_{43}(s_{(40)}), s_{(40)})\right) T_{40}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{T}_{41}(t) = T_{41}^{0} + \int_0^t \left[(b_{41})^{(8)} T_{40}(s_{(40)}) - \left((b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}(G_{43}(s_{(40)}), s_{(40)})\right) T_{41}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{T}_{42}(t) = T_{42}^{0} + \int_0^t \left[(b_{42})^{(8)} T_{41}(s_{(40)}) - \left((b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}(G_{43}(s_{(40)}), s_{(40)})\right) T_{42}(s_{(40)})\right] ds_{(40)}$<br>Where $s_{(40)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(9)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 166 A |

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| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{44})^{(9)}$ , $T_i^0 \leq (\hat{Q}_{44})^{(9)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} t}$ | |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} t}$ | |
| By<br>$\bar{G}_{44}(t) = G_{44}^0 + \int_0^t \left[(a_{44})^{(9)} G_{45}(s_{(44)}) - \left((a'_{44})^{(9)} + a''_{44})^{(9)} (T_{45}(s_{(44)}), s_{(44)})\right) G_{44}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{G}_{45}(t) = G_{45}^0 + \int_0^t \left[(a_{45})^{(9)} G_{44}(s_{(44)}) - \left((a'_{45})^{(9)} + (a''_{45})^{(9)} (T_{45}(s_{(44)}), s_{(44)})\right) G_{45}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{G}_{46}(t) = G_{46}^0 + \int_0^t \left[(a_{46})^{(9)} G_{45}(s_{(44)}) - \left((a'_{46})^{(9)} + (a''_{46})^{(9)} (T_{45}(s_{(44)}), s_{(44)})\right) G_{46}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{44}(t) = T_{44}^0 + \int_0^t \left[(b_{44})^{(9)} T_{45}(s_{(44)}) - \left((b'_{44})^{(9)} - (b''_{44})^{(9)} (G_{47}(s_{(44)}), s_{(44)})\right) T_{44}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{45}(t) = T_{45}^0 + \int_0^t \left[(b_{45})^{(9)} T_{44}(s_{(44)}) - \left((b'_{45})^{(9)} - (b''_{45})^{(9)} (G_{47}(s_{(44)}), s_{(44)})\right) T_{45}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{46}(t) = T_{46}^0 + \int_0^t \left[(b_{46})^{(9)} T_{45}(s_{(44)}) - \left((b'_{46})^{(9)} - (b''_{46})^{(9)} (G_{47}(s_{(44)}), s_{(44)})\right) T_{46}(s_{(44)})\right] ds_{(44)}$<br>Where $s_{(44)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| The operator $\mathcal{A}^{(1)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{13}(t) \leq G_{13}^0 + \int_0^t \left[(a_{13})^{(1)} \left(G_{14}^0 + (\hat{P}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} s_{(13)}}\right)\right] ds_{(13)} =$<br>$\left(1 + (a_{13})^{(1)} t\right) G_{14}^0 + \frac{(a_{13})^{(1)} (\hat{P}_{13})^{(1)}}{(\hat{M}_{13})^{(1)}} \left(e^{(\hat{M}_{13})^{(1)} t} - 1\right)$ | 167 |
| From which it follows that<br>$(G_{13}(t) - G_{13}^0) e^{-(\hat{M}_{13})^{(1)} t} \leq \frac{(a_{13})^{(1)}}{(\hat{M}_{13})^{(1)}} \left[\left((\hat{P}_{13})^{(1)} + G_{14}^0\right) e^{\left(-\frac{(\hat{P}_{13})^{(1)} + G_{14}^0}{G_{14}^0}\right)} + (\hat{P}_{13})^{(1)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 1 | 168 |
| Analogous inequalities hold also for $G_{14}$ , $G_{15}, T_{13}, T_{14}, T_{15}$ | |
| The operator $\mathcal{A}^{(2)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that | |
| $G_{16}(t) \leq G_{16}^0 + \int_0^t \left[(a_{16})^{(2)} \left(G_{17}^0 + (\hat{P}_{16})^{(6)} e^{(\hat{M}_{16})^{(2)} s_{(16)}}\right)\right] ds_{(16)} =$<br>$\left(1 + (a_{16})^{(2)} t\right) G_{17}^0 + \frac{(a_{16})^{(2)} (\hat{P}_{16})^{(2)}}{(\hat{M}_{16})^{(2)}} \left(e^{(\hat{M}_{16})^{(2)} t} - 1\right)$ | 169 |
| From which it follows that | 170 |

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| $(G_{16}(t)-G_{16}^0)e^{-(\bar{M}_{16})^{(2)}t} \le \frac{(a_{16})^{(2)}}{(\bar{M}_{16})^{(2)}}\left[((\hat{P}_{16})^{(2)}+G_{17}^0)e^{\left(-\frac{(\hat{P}_{16})^{(2)}+G_{17}^0}{G_{17}^0}\right)}+(\hat{P}_{16})^{(2)}\right]$ | |
| Analogous inequalities hold also for $G_{17}, G_{18}, T_{16}, T_{17}, T_{18}$ | |
| The operator $\mathcal{A}^{(3)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{20}(t) \le G_{20}^0 + \int_0^t \left[(a_{20})^{(3)}\left(G_{21}^0+(\hat{P}_{20})^{(3)}e^{(\bar{M}_{20})^{(3)}s_{(20)}}\right)\right] ds_{(20)} =$<br>$(1+(a_{20})^{(3)}t)G_{21}^0 + \frac{(a_{20})^{(3)}(\hat{P}_{20})^{(3)}}{(\bar{M}_{20})^{(3)}}\left(e^{(\bar{M}_{20})^{(3)}t}-1\right)$ | 171 |
| From which it follows that<br>$(G_{20}(t)-G_{20}^0)e^{-(\bar{M}_{20})^{(3)}t} \le \frac{(a_{20})^{(3)}}{(\bar{M}_{20})^{(3)}}\left[((\hat{P}_{20})^{(3)}+G_{21}^0)e^{\left(-\frac{(\hat{P}_{20})^{(3)}+G_{21}^0}{G_{21}^0}\right)}+(\hat{P}_{20})^{(3)}\right]$ | 172 |
| Analogous inequalities hold also for $G_{21}, G_{22}, T_{20}, T_{21}, T_{22}$ | |
| The operator $\mathcal{A}^{(4)}$ maps the space of functions satisfying into itself .Indeed it is obvious that<br>$G_{24}(t) \le G_{24}^0 + \int_0^t \left[(a_{24})^{(4)}\left(G_{25}^0+(\hat{P}_{24})^{(4)}e^{(\bar{M}_{24})^{(4)}s_{(24)}}\right)\right] ds_{(24)} =$<br>$(1+(a_{24})^{(4)}t)G_{25}^0 + \frac{(a_{24})^{(4)}(\hat{P}_{24})^{(4)}}{(\bar{M}_{24})^{(4)}}\left(e^{(\bar{M}_{24})^{(4)}t}-1\right)$ | 173 |
| From which it follows that<br>$(G_{24}(t)-G_{24}^0)e^{-(\bar{M}_{24})^{(4)}t} \le \frac{(a_{24})^{(4)}}{(\bar{M}_{24})^{(4)}}\left[((\hat{P}_{24})^{(4)}+G_{25}^0)e^{\left(-\frac{(\hat{P}_{24})^{(4)}+G_{25}^0}{G_{25}^0}\right)}+(\hat{P}_{24})^{(4)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 4 | 174 |
| The operator $\mathcal{A}^{(5)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{28}(t) \le G_{28}^0 + \int_0^t \left[(a_{28})^{(5)}\left(G_{29}^0+(\hat{P}_{28})^{(5)}e^{(\bar{M}_{28})^{(5)}s_{(28)}}\right)\right] ds_{(28)} =$<br>$(1+(a_{28})^{(5)}t)G_{29}^0 + \frac{(a_{28})^{(5)}(\hat{P}_{28})^{(5)}}{(\bar{M}_{28})^{(5)}}\left(e^{(\bar{M}_{28})^{(5)}t}-1\right)$ | |
| From which it follows that<br>$(G_{28}(t)-G_{28}^0)e^{-(\bar{M}_{28})^{(5)}t} \le \frac{(a_{28})^{(5)}}{(\bar{M}_{28})^{(5)}}\left[((\hat{P}_{28})^{(5)}+G_{29}^0)e^{\left(-\frac{(\hat{P}_{28})^{(5)}+G_{29}^0}{G_{29}^0}\right)}+(\hat{P}_{28})^{(5)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 5 | 175 |
| The operator $\mathcal{A}^{(6)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious | 176 |

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| that<br>$G_{32}(t) \le G_{32}^0 + \int_0^t \left[(a_{32})^{(6)}\left(G_{33}^0 + (\hat{P}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} s_{(32)}}\right)\right] ds_{(32)} =$<br>$\left(1 + (a_{32})^{(6)} t\right) G_{33}^0 + \frac{(a_{32})^{(6)} (\hat{P}_{32})^{(6)}}{(\hat{M}_{32})^{(6)}}\left(e^{(\hat{M}_{32})^{(6)} t} - 1\right)$ | |
| From which it follows that<br>$(G_{32}(t) - G_{32}^0) e^{-(\hat{M}_{32})^{(6)} t} \le \frac{(a_{32})^{(6)}}{(\hat{M}_{32})^{(6)}}\left[\left((\hat{P}_{32})^{(6)} + G_{33}^0\right) e^{\left(-\frac{(\hat{P}_{32})^{(6)} + G_{33}^0}{G_{33}^0}\right)} + (\hat{P}_{32})^{(6)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 6<br>Analogous inequalities hold also for $G_{25}, G_{26}, T_{24}, T_{25}, T_{26}$ | 177 |
| (g) The operator $\mathcal{A}^{(7)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that<br>$G_{36}(t) \le G_{36}^0 + \int_0^t \left[(a_{36})^{(7)}\left(G_{37}^0 + (\hat{P}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} s_{(36)}}\right)\right] ds_{(36)} =$<br>$\left(1 + (a_{36})^{(7)} t\right) G_{37}^0 + \frac{(a_{36})^{(7)} (\hat{P}_{36})^{(7)}}{(\hat{M}_{36})^{(7)}}\left(e^{(\hat{M}_{36})^{(7)} t} - 1\right)$ | 178 |
| From which it follows that<br>$(G_{36}(t) - G_{36}^0) e^{-(\hat{M}_{36})^{(7)} t} \le \frac{(a_{36})^{(7)}}{(\hat{M}_{36})^{(7)}}\left[\left((\hat{P}_{36})^{(7)} + G_{37}^0\right) e^{\left(-\frac{(\hat{P}_{36})^{(7)} + G_{37}^0}{G_{37}^0}\right)} + (\hat{P}_{36})^{(7)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 7 | |
| The operator $\mathcal{A}^{(8)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that | |
| $G_{40}(t) \le G_{40}^0 + \int_0^t \left[(a_{40})^{(8)}\left(G_{41}^0 + (\hat{P}_{40})^{(8)} e^{(\hat{M}_{40})^{(8)} s_{(40)}}\right)\right] ds_{(40)} =$<br>$\left(1 + (a_{40})^{(8)} t\right) G_{41}^0 + \frac{(a_{40})^{(8)} (\hat{P}_{40})^{(8)}}{(\hat{M}_{40})^{(8)}}\left(e^{(\hat{M}_{40})^{(8)} t} - 1\right)$ | 180 |
| From which it follows that<br>$(G_{40}(t) - G_{40}^0) e^{-(\hat{M}_{40})^{(8)} t} \le \frac{(a_{40})^{(8)}}{(\hat{M}_{40})^{(8)}}\left[\left((\hat{P}_{40})^{(8)} + G_{41}^0\right) e^{\left(-\frac{(\hat{P}_{40})^{(8)} + G_{41}^0}{G_{41}^0}\right)} + (\hat{P}_{40})^{(8)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 8<br>Analogous inequalities hold also for $G_{41}, G_{42}, T_{40}, T_{41}, T_{42}$ | 181 |
| The operator $\mathcal{A}^{(9)}$ maps the space of functions satisfying 34,35,36 into itself .Indeed it is obvious that<br>$G_{44}(t) \le G_{44}^0 + \int_0^t \left[(a_{44})^{(9)}\left(G_{45}^0 + (\hat{P}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} s_{(44)}}\right)\right] ds_{(44)} =$ | |

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| $\left(1+(a_{44})^{(9)}t\right)G_{45}^{0}+\frac{(a_{44})^{(9)}(\hat{P}_{44})^{(9)}}{(\hat{M}_{44})^{(9)}}\left(e^{(\hat{M}_{44})^{(9)}t}-1\right)$ | |
| From which it follows that<br>$(G_{44}(t)-G_{44}^{0})e^{-(\hat{M}_{44})^{(9)}t}\leq\frac{(a_{44})^{(9)}}{(\hat{M}_{44})^{(9)}}\left[\left((\hat{P}_{44})^{(9)}+G_{45}^{0}\right)e^{\left(-\frac{(\hat{P}_{44})^{(9)}+G_{45}^{0}}{G_{45}^{0}}\right)}+(\hat{P}_{44})^{(9)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 9<br>Analogous inequalities hold also for $G_{45}, G_{46}, T_{44}, T_{45}, T_{46}$ | |
| It is now sufficient to take $\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}}<1$ and to choose<br>$(\hat{P}_{13})^{(1)}$ and $(\hat{Q}_{13})^{(1)}$ large to have | 182 |
| $\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}\left[(\hat{P}_{13})^{(1)}+\left((\hat{P}_{13})^{(1)}+G_j^0\right)e^{-\left(\frac{(\hat{P}_{13})^{(1)}+G_j^0}{G_j^0}\right)}\right]\leq(\hat{P}_{13})^{(1)}$ | 183 |
| $\frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}}\left[\left((\hat{Q}_{13})^{(1)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{13})^{(1)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{13})^{(1)}\right]\leq(\hat{Q}_{13})^{(1)}$ | 184 |
| In order that the operator $\mathcal{A}^{(1)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(1)}$ is a contraction with respect to the metric<br>$d\left(\left(G^{(1)},T^{(1)}\right),\left(G^{(2)},T^{(2)}\right)\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\hat{M}_{13})^{(1)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\hat{M}_{13})^{(1)}t}\}$ | 185 |
| Indeed if we denote<br>**Definition of** $\tilde{G},\tilde{T}:\left(\tilde{G},\tilde{T}\right)=\mathcal{A}^{(1)}(G,T)$<br>It results<br>$\left\|\tilde{G}_{13}^{(1)}-\tilde{G}_i^{(2)}\right\|\leq\int_0^t(a_{13})^{(1)}\left\|G_{14}^{(1)}-G_{14}^{(2)}\right\|e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{(\hat{M}_{13})^{(1)}s_{(13)}}ds_{(13)}+$<br>$\int_0^t\{(a_{13}')^{(1)}\left\|G_{13}^{(1)}-G_{13}^{(2)}\right\|e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{-(\hat{M}_{13})^{(1)}s_{(13)}}+$<br>$(a_{13}'')^{(1)}\left(T_{14}^{(1)},s_{(13)}\right)\left\|G_{13}^{(1)}-G_{13}^{(2)}\right\|e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{(\hat{M}_{13})^{(1)}s_{(13)}}+$<br>$G_{13}^{(2)}\left\\|(a_{13}'')^{(1)}\left(T_{14}^{(1)},s_{(13)}\right)-(a_{13}'')^{(1)}\left(T_{14}^{(2)},s_{(13)}\right)\right\\|\ e^{-(\hat{M}_{13})^{(1)}s_{(13)}}e^{(\hat{M}_{13})^{(1)}s_{(13)}}\}ds_{(13)}$ | |

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| Where $s_{(13)}$ represents integrand that is integrated over the interval $[0, t]$<br><br>From the hypotheses it follows | |
| $\left\|G^{(1)} - G^{(2)}\right\| e^{-(\widehat{M}_{13})^{(1)}t} \leq$<br>$\frac{1}{(\widehat{M}_{13})^{(1)}}\left((a_{13})^{(1)} + (a'_{13})^{(1)} + (\widehat{A}_{13})^{(1)} + (\widehat{P}_{13})^{(1)}(\widehat{k}_{13})^{(1)}\right) d\left(\left(G^{(1)}, T^{(1)};\ G^{(2)}, T^{(2)}\right)\right)$<br><br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 186 |
| **Remark 1:** The fact that we supposed $(a''_{13})^{(1)}$ and $(b''_{13})^{(1)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{13})^{(1)} e^{(\widehat{M}_{13})^{(1)}t}$ *and* $(\widehat{Q}_{13})^{(1)} e^{(\widehat{M}_{13})^{(1)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a''_i)^{(1)}$ and $(b''_i)^{(1)}, i = 13,14,15$ depend only on $T_{14}$ and respectively on $G(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 2:** There does not exist any $t$ where $G_i(t) = 0$ *and* $T_i(t) = 0$<br><br>From 19 to 24 it results<br><br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a'_i)^{(1)} - (a''_i)^{(1)}(T_{14}(s_{(13)}), s_{(13)})\} ds_{(13)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{\left(-(b'_i)^{(1)}t\right)} > 0$ for t > 0 | |
| **Definition of** $\left((\widehat{M}_{13})^{(1)}\right)_1, \left((\widehat{M}_{13})^{(1)}\right)_2$ *and* $\left((\widehat{M}_{13})^{(1)}\right)_3$ :<br><br>**Remark 3**: if $G_{13}$ is bounded, the same property have also $G_{14}$ *and* $G_{15}$ . indeed if<br><br>$G_{13} < (\widehat{M}_{13})^{(1)}$ it follows $\frac{dG_{14}}{dt} \leq \left((\widehat{M}_{13})^{(1)}\right)_1 - (a'_{14})^{(1)} G_{14}$ and by integrating<br><br>$G_{14} \leq \left((\widehat{M}_{13})^{(1)}\right)_2 = G_{14}^0 + 2(a_{14})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_1 / (a'_{14})^{(1)}$<br><br>In the same way , one can obtain<br><br>$G_{15} \leq \left((\widehat{M}_{13})^{(1)}\right)_3 = G_{15}^0 + 2(a_{15})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_2 / (a'_{15})^{(1)}$<br><br>If $G_{14}$ *or* $G_{15}$ is bounded, the same property follows for $G_{13}$ , $G_{15}$ and $G_{13}$ , $G_{14}$ respectively. | 187 |
| **Remark 4:** If $G_{13}$ *is* bounded, from below, the same property holds for $G_{14}$ *and* $G_{15}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{14}$ is bounded from below. | 188 |
| **Remark 5:** If $T_{13}$ is bounded from below and $\lim_{t\to\infty}((b''_i)^{(1)}(G(t), t)) = (b'_{14})^{(1)}$ then $T_{14} \to \infty$.<br><br>**Definition of** $(m)^{(1)}$ and $\varepsilon_1$ : | 189 |

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| Indeed let $t_1$ be so that for $t > t_1$<br>$(b_{14})^{(1)} - (b_i^{''})^{(1)}(G(t),t) < \varepsilon_1, T_{13}(t) > (m)^{(1)}$ | |
| Then $\frac{dT_{14}}{dt} \geq (a_{14})^{(1)}(m)^{(1)} - \varepsilon_1 T_{14}$ which leads to<br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{\varepsilon_1}\right)(1 - e^{-\varepsilon_1 t}) + T_{14}^0 e^{-\varepsilon_1 t}$ If we take t such that $e^{-\varepsilon_1 t} = \frac{1}{2}$ it results<br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{2}\right), \quad t = log\frac{2}{\varepsilon_1}$ By taking now $\varepsilon_1$ sufficiently small one sees that $T_{14}$ is unbounded.<br>The same property holds for $T_{15}$ if $\lim_{t\to\infty}(b_{15}^{''})^{(1)}(G(t),t) = (b_{15}^{'})^{(1)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |
| It is now sufficient to take $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}$ , $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ and to choose<br>$(\hat{P}_{16})^{(2)}$ $and$ $(\hat{Q}_{16})^{(2)}$ large to have | 190 |
| $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[(\hat{P}_{16})^{(2)} + \left((\hat{P}_{16})^{(2)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{16})^{(2)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{16})^{(2)}$ | 191 |
| $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[\left((\hat{Q}_{16})^{(2)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{16})^{(2)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{16})^{(2)}\right] \leq (\hat{Q}_{16})^{(2)}$ | 192 |
| In order that the operator $\mathcal{A}^{(2)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 193 |
| The operator $\mathcal{A}^{(2)}$ is a contraction with respect to the metric<br>$d\left(((G_{19})^{(1)},(T_{19})^{(1)}),((G_{19})^{(2)},(T_{19})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\hat{M}_{16})^{(2)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\hat{M}_{16})^{(2)}t}\}$ | 194 |
| Indeed if we denote<br>**Definition of** $\widetilde{G_{19}}, \widetilde{T_{19}}$ : $\left(\widetilde{G_{19}}, \widetilde{T_{19}}\right) = \mathcal{A}^{(2)}(G_{19}, T_{19})$ | 195 |

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| It results<br>$\left\|\tilde{G}_{16}^{(1)} - \tilde{G}_{i}^{(2)}\right\| \leq \int_0^t (a_{16})^{(2)} \left\|G_{17}^{(1)} - G_{17}^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{(\widehat{M}_{16})^{(2)} s_{(16)}}\, ds_{(16)} +$<br>$\int_0^t \{(a_{16}^{'})^{(2)} \left\|G_{16}^{(1)} - G_{16}^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} +$<br>$(a_{16}^{''})^{(2)} \left(T_{17}^{(1)}, s_{(16)}\right) \left\|G_{16}^{(1)} - G_{16}^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{(\widehat{M}_{16})^{(2)} s_{(16)}} +$<br>$G_{16}^{(2)} \left\|(a_{16}^{''})^{(2)} \left(T_{17}^{(1)}, s_{(16)}\right) - (a_{16}^{''})^{(2)} \left(T_{17}^{(2)}, s_{(16)}\right)\right\| \; e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{(\widehat{M}_{16})^{(2)} s_{(16)}} \} ds_{(16)}$ | 196 |
| Where $s_{(16)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | 197 |
| $\left\|(G_{19})^{(1)} - (G_{19})^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} t} \leq$<br>$\frac{1}{(\widehat{M}_{16})^{(2)}} \left((a_{16})^{(2)} + (a_{16}^{'})^{(2)} + (\widehat{A}_{16})^{(2)} + (\widehat{P}_{16})^{(2)} (\widehat{k}_{16})^{(2)}\right) d\left(\left((G_{19})^{(1)}, (T_{19})^{(1)}; (G_{19})^{(2)}, (T_{19})^{(2)}\right)\right)$ | |
| And analogous inequalities for $G_i$ and $T_i$. Taking into account the hypothesis the result follows | 198 |
| **Remark 6:** The fact that we supposed $(a_{16}^{''})^{(2)}$ and $(b_{16}^{''})^{(2)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{16})^{(2)} e^{(\widehat{M}_{16})^{(2)} t}$ and $(\widehat{Q}_{16})^{(2)} e^{(\widehat{M}_{16})^{(2)} t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}, i = 16,17,18$ depend only on $T_{17}$ and respectively on $(G_{19})$(and not on t) and hypothesis can replaced by a usual Lipschitz condition. | 199 |
| **Remark 7:** There does not exist any t where $G_i(t) = 0$ and $T_i(t) = 0$<br>it results<br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(2)} - (a_i^{''})^{(2)} (T_{17}(s_{(16)}), s_{(16)})\} ds_{(16)}\right]} \geq 0$<br>$T_i(t) \geq T_i^0 e^{(-(b_i^{'})^{(2)} t)} > 0 \quad \text{for } t > 0$ | 200 |
| **Definition of** $\left((\widehat{M}_{16})^{(2)}\right)_1$, $\left((\widehat{M}_{16})^{(2)}\right)_2$ and $\left((\widehat{M}_{16})^{(2)}\right)_3$ :<br>**Remark 8:** if $G_{16}$ is bounded, the same property have also $G_{17}$ and $G_{18}$ . indeed if<br>$G_{16} < (\widehat{M}_{16})^{(2)}$ it follows $\frac{dG_{17}}{dt} \leq \left((\widehat{M}_{16})^{(2)}\right)_1 - (a_{17}^{'})^{(2)} G_{17}$ and by integrating<br>$G_{17} \leq \left((\widehat{M}_{16})^{(2)}\right)_2 = G_{17}^0 + 2(a_{17})^{(2)} \left((\widehat{M}_{16})^{(2)}\right)_1 / (a_{17}^{'})^{(2)}$<br>In the same way , one can obtain | 201 |

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| $G_{18} \leq \left((\widehat{M}_{16})^{(2)}\right)_3 = G_{18}^0 + 2(a_{18})^{(2)}\left((\widehat{M}_{16})^{(2)}\right)_2/(a_{18}')^{(2)}$<br><br>If $G_{17}$ or $G_{18}$ is bounded, the same property follows for $G_{16}$, $G_{18}$ and $G_{16}$, $G_{17}$ respectively. | |
| **Remark 9**: If $G_{16}$ is bounded, from below, the same property holds for $G_{17}$ and $G_{18}$. The proof is analogous with the preceding one. An analogous property is true if $G_{17}$ is bounded from below. | 202 |
| **Remark 10:** If $T_{16}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(2)}((G_{19})(t),t)) = (b_{17}')^{(2)}$ then $T_{17} \to \infty$.<br><br>**Definition of** $(m)^{(2)}$ and $\varepsilon_2$ :<br><br>Indeed let $t_2$ be so that for $t > t_2$<br><br>$(b_{17})^{(2)} - (b_i'')^{(2)}((G_{19})(t),t) < \varepsilon_2, T_{16}(t) > (m)^{(2)}$ | 203 |
| Then $\frac{dT_{17}}{dt} \geq (a_{17})^{(2)}(m)^{(2)} - \varepsilon_2 T_{17}$ which leads to<br><br>$T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{\varepsilon_2}\right)(1 - e^{-\varepsilon_2 t}) + T_{17}^0 e^{-\varepsilon_2 t}$ If we take t such that $e^{-\varepsilon_2 t} = \frac{1}{2}$ it results | 204 |
| $T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{2}\right)$, $t = \log\frac{2}{\varepsilon_2}$ By taking now $\varepsilon_2$ sufficiently small one sees that $T_{17}$ is unbounded.<br>The same property holds for $T_{18}$ if $\lim_{t\to\infty}(b_{18}'')^{(2)}\left((G_{19})(t),t\right) = (b_{18}')^{(2)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 205 |
| It is now sufficient to take $\frac{(a_i)^{(3)}}{(\bar{M}_{20})^{(3)}}$, $\frac{(b_i)^{(3)}}{(\bar{M}_{20})^{(3)}} < 1$ and to choose<br><br>$(\hat{P}_{20})^{(3)}$ and $(\hat{Q}_{20})^{(3)}$ large to have | 207 |
| $\frac{(a_i)^{(3)}}{(\bar{M}_{20})^{(3)}}\left[(\hat{P}_{20})^{(3)} + \left((\hat{P}_{20})^{(3)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{20})^{(3)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{20})^{(3)}$ | 208 |
| $\frac{(b_i)^{(3)}}{(\bar{M}_{20})^{(3)}}\left[\left((\hat{Q}_{20})^{(3)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{20})^{(3)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{20})^{(3)}\right] \leq (\hat{Q}_{20})^{(3)}$ | 209 |
| In order that the operator $\mathcal{A}^{(3)}$ transforms the space of sextuples of functions $G_i$, $T_i$ satisfying Equations into itself | 210 |
| The operator $\mathcal{A}^{(3)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{23})^{(1)},(T_{23})^{(1)}),((G_{23})^{(2)},(T_{23})^{(2)})\right) =$<br><br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t) - G_i^{(2)}(t)\right|e^{-(\bar{M}_{20})^{(3)}t}, \max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t) - T_i^{(2)}(t)\right|e^{-(\bar{M}_{20})^{(3)}t}\}$ | 211 |

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| Indeed if we denote<br><br>**Definition of** $\widetilde{G_{23}}, \widetilde{T_{23}} :\left( \widetilde{(G_{23})}, \widetilde{(T_{23})} \right) = \mathcal{A}^{(3)}\left((G_{23}), (T_{23})\right)$ | 212 |
| It results<br><br>$\left\|\tilde{G}_{20}^{(1)} - \tilde{G}_{i}^{(2)}\right\| \leq \int_0^t (a_{20})^{(3)} \left\|G_{21}^{(1)} - G_{21}^{(2)}\right\| e^{-(\widehat{M}_{20})^{(3)} s_{(20)}} e^{(\widehat{M}_{20})^{(3)} s_{(20)}}\, ds_{(20)} +$<br><br>$\int_0^t \{(a_{20}^{'})^{(3)} \left\|G_{20}^{(1)} - G_{20}^{(2)}\right\| e^{-(\widehat{M}_{20})^{(3)} s_{(20)}} e^{-(\widehat{M}_{20})^{(3)} s_{(20)}} +$<br><br>$(a_{20}^{''})^{(3)}\left(T_{21}^{(1)}, s_{(20)}\right)\left\|G_{20}^{(1)} - G_{20}^{(2)}\right\| e^{-(\widehat{M}_{20})^{(3)} s_{(20)}} e^{(\widehat{M}_{20})^{(3)} s_{(20)}} +$<br><br>$G_{20}^{(2)} \left\|(a_{20}^{''})^{(3)}\left(T_{21}^{(1)}, s_{(20)}\right) - (a_{20}^{''})^{(3)}\left(T_{21}^{(2)}, s_{(20)}\right)\right\| \; e^{-(\widehat{M}_{20})^{(3)} s_{(20)}} e^{(\widehat{M}_{20})^{(3)} s_{(20)}}\} ds_{(20)}$<br><br>Where $s_{(20)}$ represents integrand that is integrated over the interval $[0, t]$<br><br>From the hypotheses it follows | 213 |
| $\left\|G_{23}{}^{(1)} - G_{23}{}^{(2)}\right\| e^{-(\widehat{M}_{20})^{(3)} t} \leq$<br>$\frac{1}{(\widehat{M}_{20})^{(3)}}\left((a_{20})^{(3)} + (a_{20}^{'})^{(3)} + (\widehat{A}_{20})^{(3)} + (\widehat{P}_{20})^{(3)} (\widehat{k}_{20})^{(3)}\right) d\left(\left((G_{23})^{(1)}, (T_{23})^{(1)};\ (G_{23})^{(2)}, (T_{23})^{(2)}\right)\right)$<br><br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 214 |
| **Remark 11:** The fact that we supposed $(a_{20}^{''})^{(3)}$ and $(b_{20}^{''})^{(3)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{20})^{(3)} e^{(\widehat{M}_{20})^{(3)} t}$ *and* $(\widehat{Q}_{20})^{(3)} e^{(\widehat{M}_{20})^{(3)} t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(3)}$ and $(b_i^{''})^{(3)}, i = 20,21,22$ depend only on $T_{21}$ and respectively on $(G_{23})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 215 |
| **Remark 12:** There does not exist any $t$ where $G_i\,(t) = 0$ *and* $T_i\,(t) = 0$<br><br>it results<br><br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(3)} - (a_i^{''})^{(3)} (T_{21}(s_{(20)}), s_{(20)})\} ds_{(20)}\right]} \geq 0$<br><br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(3)} t\right)} > 0$ for t > 0 | 216 |
| **Definition of** $\left((\widehat{M}_{20})^{(3)}\right)_1, \left((\widehat{M}_{20})^{(3)}\right)_2$ *and* $\left((\widehat{M}_{20})^{(3)}\right)_3$ :<br><br>**Remark 13:** if $G_{20}$ is bounded, the same property have also $G_{21}$ *and* $G_{22}$ . indeed if<br><br>$G_{20} < (\widehat{M}_{20})^{(3)}$ it follows $\frac{dG_{21}}{dt} \leq \left((\widehat{M}_{20})^{(3)}\right)_1 - (a_{21}^{'})^{(3)} G_{21}$ and by integrating | 217 |

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| $G_{21} \leq \left((\widehat{M}_{20})^{(3)}\right)_2 = G_{21}^0 + 2(a_{21})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_1/(a_{21}')^{(3)}$<br><br>In the same way , one can obtain<br><br>$G_{22} \leq \left((\widehat{M}_{20})^{(3)}\right)_3 = G_{22}^0 + 2(a_{22})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_2/(a_{22}')^{(3)}$<br><br>If $G_{21}$ *or* $G_{22}$ is bounded, the same property follows for $G_{20}$ , $G_{22}$ and $G_{20}$ , $G_{21}$ respectively. | |
| **Remark 14:** If $G_{20}$ *is* bounded, from below, the same property holds for $G_{21}$ *and* $G_{22}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{21}$ is bounded from below. | 218 |
| **Remark 15:** If $T_{20}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(3)}\left((G_{23})(t),t\right) = (b_{21}^{'})^{(3)}$ then $T_{21} \to \infty$.<br><br>**Definition of** $(m)^{(3)}$ and $\varepsilon_3$ :<br><br>Indeed let $t_3$ be so that for $t > t_3$<br><br>$(b_{21})^{(3)} - (b_i^{''})^{(3)}\left((G_{23})(t),t\right) < \varepsilon_3, T_{20}(t) > (m)^{(3)}$ | 219 |
| Then $\frac{dT_{21}}{dt} \geq (a_{21})^{(3)}(m)^{(3)} - \varepsilon_3 T_{21}$ which leads to<br><br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{\varepsilon_3}\right)(1 - e^{-\varepsilon_3 t}) + T_{21}^0 e^{-\varepsilon_3 t}$ If we take t such that $e^{-\varepsilon_3 t} = \frac{1}{2}$ it results<br><br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{2}\right)$, $t = log\frac{2}{\varepsilon_3}$ By taking now $\varepsilon_3$ sufficiently small one sees that $T_{21}$ is unbounded.<br>The same property holds for $T_{22}$ if $\lim_{t\to\infty}(b_{22}^{''})^{(3)}\left((G_{23})(t),t\right) = (b_{22}^{'})^{(3)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 220 |
| It is now sufficient to take $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}$ , $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}} < 1$ and to choose<br><br>$(\widehat{P}_{24})^{(4)}$ and $(\widehat{Q}_{24})^{(4)}$ large to have | 221 |
| $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[(\widehat{P}_{24})^{(4)} + \left((\widehat{P}_{24})^{(4)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{24})^{(4)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{24})^{(4)}$ | 222 |
| $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[\left((\widehat{Q}_{24})^{(4)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{24})^{(4)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{24})^{(4)}\right] \leq (\widehat{Q}_{24})^{(4)}$ | 223 |
| In order that the operator $\mathcal{A}^{(4)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 224 |
| The operator $\mathcal{A}^{(4)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{27})^{(1)},(T_{27})^{(1)}),((G_{27})^{(2)},(T_{27})^{(2)})\right) =$ | 225 |

$$\sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t)-G_i^{(2)}(t)\right|e^{-(\widehat{M}_{24})^{(4)}t},\max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t)-T_i^{(2)}(t)\right|e^{-(\widehat{M}_{24})^{(4)}t}\}$$

Indeed if we denote

**Definition of** $\widetilde{(G_{27})},\widetilde{(T_{27})}$ : $\left(\widetilde{(G_{27})},\widetilde{(T_{27})}\right)=\mathcal{A}^{(4)}((G_{27}),(T_{27}))$

It results

$$\left|\tilde{G}_{24}^{(1)}-\tilde{G}_i^{(2)}\right|\le\int_0^t(a_{24})^{(4)}\left|G_{25}^{(1)}-G_{25}^{(2)}\right|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}\,ds_{(24)}+$$

$$\int_0^t\{(a_{24}^{'})^{(4)}\left|G_{24}^{(1)}-G_{24}^{(2)}\right|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}+$$

$$(a_{24}^{''})^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)\left|G_{24}^{(1)}-G_{24}^{(2)}\right|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}+$$

$$G_{24}^{(2)}\left|(a_{24}^{''})^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)-(a_{24}^{''})^{(4)}\left(T_{25}^{(2)},s_{(24)}\right)\right|\ e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}\}ds_{(24)}$$

Where $s_{(24)}$ represents integrand that is integrated over the interval $[0,t]$

From the hypotheses on Equations it follows

$$\left|(G_{27})^{(1)}-(G_{27})^{(2)}\right|e^{-(\widehat{M}_{24})^{(4)}t}\le$$
$$\frac{1}{(\widehat{M}_{24})^{(4)}}\left((a_{24})^{(4)}+(a_{24}^{'})^{(4)}+(\widehat{A}_{24})^{(4)}+(\widehat{P}_{24})^{(4)}(\widehat{k}_{24})^{(4)}\right)d\left(((G_{27})^{(1)},(T_{27})^{(1)};\ (G_{27})^{(2)},(T_{27})^{(2)})\right)$$

(226)

And analogous inequalities for $G_i\ and\ T_i$. Taking into account the hypothesis the result follows

**Remark 16:** The fact that we supposed $(a_{24}^{''})^{(4)}$ and $(b_{24}^{''})^{(4)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}\ and\ (\widehat{Q}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}$ respectively of $\mathbb{R}_+$.

(227)

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(4)}$ and $(b_i^{''})^{(4)}, i=24,25,26$ depend only on $T_{25}$ and respectively on $(G_{27})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition.

**Remark 17:** There does not exist any $t$ where $G_i\ (t)=0\ and\ T_i\ (t)=0$

(228)

it results

$$G_i\ (t)\ge G_i^0e^{\left[-\int_0^t\{(a_i^{'})^{(4)}-(a_i^{''})^{(4)}(T_{25}(s_{(24)}),s_{(24)})\}ds_{(24)}\right]}\ge 0$$

$$T_i\ (t)\ge T_i^0e^{\left(-(b_i^{'})^{(4)}t\right)}>0\quad \text{for t} > 0$$

**Definition of** $\left((\widehat{M}_{24})^{(4)}\right)_1,\left((\widehat{M}_{24})^{(4)}\right)_2\ and\ \left((\widehat{M}_{24})^{(4)}\right)_3$ :

(229)

**Remark 18:** if $G_{24}$ is bounded, the same property have also $G_{25}\ and\ G_{26}$ . indeed if

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| $G_{24} < (\widehat{M}_{24})^{(4)}$ it follows $\frac{dG_{25}}{dt} \leq \left((\widehat{M}_{24})^{(4)}\right)_1 - (a'_{25})^{(4)} G_{25}$ and by integrating<br>$G_{25} \leq \left((\widehat{M}_{24})^{(4)}\right)_2 = G_{25}^0 + 2(a_{25})^{(4)}\left((\widehat{M}_{24})^{(4)}\right)_1/(a'_{25})^{(4)}$<br>In the same way , one can obtain<br>$G_{26} \leq \left((\widehat{M}_{24})^{(4)}\right)_3 = G_{26}^0 + 2(a_{26})^{(4)}\left((\widehat{M}_{24})^{(4)}\right)_2/(a'_{26})^{(4)}$<br>If $G_{25}$ *or* $G_{26}$ is bounded, the same property follows for $G_{24}$ , $G_{26}$ and $G_{24}$ , $G_{25}$ respectively. | |
| **Remark 19:** If $G_{24}$ *is* bounded, from below, the same property holds for $G_{25}$ *and* $G_{26}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{25}$ is bounded from below. | 230 |
| **Remark 20:** If $T_{24}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(4)}((G_{27})(t),t)) = (b'_{25})^{(4)}$ then $T_{25} \to \infty$.<br>**Definition of** $(m)^{(4)}$ and $\varepsilon_4$ :<br>Indeed let $t_4$ be so that for $t > t_4$<br>$(b_{25})^{(4)} - (b_i'')^{(4)}((G_{27})(t),t) < \varepsilon_4, T_{24}(t) > (m)^{(4)}$ | 231 |
| Then $\frac{dT_{25}}{dt} \geq (a_{25})^{(4)}(m)^{(4)} - \varepsilon_4 T_{25}$ which leads to<br>$T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{\varepsilon_4}\right)(1 - e^{-\varepsilon_4 t}) + T_{25}^0 e^{-\varepsilon_4 t}$ If we take t such that $e^{-\varepsilon_4 t} = \frac{1}{2}$ it results<br>$T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{2}\right)$, $t = log\frac{2}{\varepsilon_4}$ By taking now $\varepsilon_4$ sufficiently small one sees that $T_{25}$ is unbounded.<br>The same property holds for $T_{26}$ if $\lim_{t\to\infty}(b''_{26})^{(4)}\left((G_{27})(t),t\right) = (b'_{26})^{(4)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 42<br>Analogous inequalities hold also for $G_{29}$ , $G_{30}, T_{28}, T_{29}, T_{30}$ | 232 |
| It is now sufficient to take $\frac{(a_i)^{(5)}}{(\widehat{M}_{28})^{(5)}}$ , $\frac{(b_i)^{(5)}}{(\widehat{M}_{28})^{(5)}} < 1$ and to choose<br>$(\widehat{P}_{28})^{(5)}$ and $(\widehat{Q}_{28})^{(5)}$ large to have | 233 |
| $\frac{(a_i)^{(5)}}{(\widehat{M}_{28})^{(5)}}\left[(\widehat{P}_{28})^{(5)} + \left((\widehat{P}_{28})^{(5)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{28})^{(5)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{28})^{(5)}$ | 234 |
| $\frac{(b_i)^{(5)}}{(\widehat{M}_{28})^{(5)}}\left[\left((\widehat{Q}_{28})^{(5)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{28})^{(5)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{28})^{(5)}\right] \leq (\widehat{Q}_{28})^{(5)}$ | 235 |
| In order that the operator $\mathcal{A}^{(5)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying | |

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| Equations into itself | |
| The operator $\mathcal{A}^{(5)}$ is a contraction with respect to the metric<br>$d\left(((G_{31})^{(1)},(T_{31})^{(1)}),((G_{31})^{(2)},(T_{31})^{(2)})\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{28})^{(5)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{28})^{(5)}t}\}$<br>Indeed if we denote<br>**Definition of** $(\widetilde{G_{31}}),(\widetilde{T_{31}})$ : $\left((\widetilde{G_{31}}),(\widetilde{T_{31}})\right)=\mathcal{A}^{(5)}((G_{31}),(T_{31}))$<br>It results<br>$\left\|\tilde{G}_{28}^{(1)}-\tilde{G}_i^{(2)}\right\|\le\int_0^t(a_{28})^{(5)}\left\|G_{29}^{(1)}-G_{29}^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{(\widehat{M}_{28})^{(5)}s_{(28)}}\,ds_{(28)}+$<br>$\int_0^t\{(a_{28}^{'})^{(5)}\left\|G_{28}^{(1)}-G_{28}^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}+$<br>$(a_{28}^{''})^{(5)}\left(T_{29}^{(1)},s_{(28)}\right)\left\|G_{28}^{(1)}-G_{28}^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{(\widehat{M}_{28})^{(5)}s_{(28)}}+$<br>$G_{28}^{(2)}\left\|(a_{28}^{''})^{(5)}\left(T_{29}^{(1)},s_{(28)}\right)-(a_{28}^{''})^{(5)}\left(T_{29}^{(2)},s_{(28)}\right)\right\|\ e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{(\widehat{M}_{28})^{(5)}s_{(28)}}\}ds_{(28)}$<br>Where $s_{(28)}$ represents integrand that is integrated over the interval [0, t]<br>From the hypotheses on it follows | 236 |
| $\left\|(G_{31})^{(1)}-(G_{31})^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}t}\le$<br>$\frac{1}{(\widehat{M}_{28})^{(5)}}\left((a_{28})^{(5)}+(a_{28}^{'})^{(5)}+(\widehat{A}_{28})^{(5)}+(\widehat{P}_{28})^{(5)}(\widehat{k}_{28})^{(5)}\right)d\left(((G_{31})^{(1)},(T_{31})^{(1)};\ (G_{31})^{(2)},(T_{31})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 237 |
| **Remark 21:** The fact that we supposed $(a_{28}^{''})^{(5)}$ and $(b_{28}^{''})^{(5)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{28})^{(5)}e^{(\widehat{M}_{28})^{(5)}t}$ *and* $(\widehat{Q}_{28})^{(5)}e^{(\widehat{M}_{28})^{(5)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(5)}$ and $(b_i^{''})^{(5)}, i=28,29,30$ depend only on $T_{29}$ and respectively on $(G_{31})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 238 |
| **Remark 22:** There does not exist any $t$ where $G_i(t)=0$ *and* $T_i(t)=0$<br>it results<br>$G_i(t)\ge G_i^0e^{\left[-\int_0^t\{(a_i^{'})^{(5)}-(a_i^{''})^{(5)}(T_{29}(s_{(28)}),s_{(28)})\}ds_{(28)}\right]}\ge 0$ | 239 |

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| $T_i(t) \geq T_i^0 e^{(-(b_i')^{(5)}t)} > 0$ for t > 0 | |
| **Definition of** $((\widehat{M}_{28})^{(5)})_1, ((\widehat{M}_{28})^{(5)})_2$ *and* $((\widehat{M}_{28})^{(5)})_3$ :<br><br>**Remark 23:** if $G_{28}$ is bounded, the same property have also $G_{29}$ *and* $G_{30}$ . indeed if<br><br>$G_{28} < (\widehat{M}_{28})^{(5)}$ it follows $\frac{dG_{29}}{dt} \leq ((\widehat{M}_{28})^{(5)})_1 - (a_{29}')^{(5)} G_{29}$ and by integrating<br><br>$G_{29} \leq ((\widehat{M}_{28})^{(5)})_2 = G_{29}^0 + 2(a_{29})^{(5)}((\widehat{M}_{28})^{(5)})_1/(a_{29}')^{(5)}$<br><br>In the same way , one can obtain<br><br>$G_{30} \leq ((\widehat{M}_{28})^{(5)})_3 = G_{30}^0 + 2(a_{30})^{(5)}((\widehat{M}_{28})^{(5)})_2/(a_{30}')^{(5)}$<br><br>If $G_{29}$ *or* $G_{30}$ is bounded, the same property follows for $G_{28}$ , $G_{30}$ and $G_{28}$ , $G_{29}$ respectively. | 240 |
| **Remark 24:** If $G_{28}$ *is* bounded, from below, the same property holds for $G_{29}$ *and* $G_{30}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{29}$ is bounded from below. | 241 |
| **Remark 25:** If $T_{28}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(5)}((G_{31})(t),t)) = (b_{29}')^{(5)}$ then $T_{29} \to \infty$.<br><br>**Definition of** $(m)^{(5)}$ and $\varepsilon_5$ :<br><br>Indeed let $t_5$ be so that for $t > t_5$<br><br>$(b_{29})^{(5)} - (b_i'')^{(5)}((G_{31})(t),t) < \varepsilon_5, T_{28}(t) > (m)^{(5)}$ | 242 |
| Then $\frac{dT_{29}}{dt} \geq (a_{29})^{(5)}(m)^{(5)} - \varepsilon_5 T_{29}$ which leads to<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{\varepsilon_5}\right)(1 - e^{-\varepsilon_5 t}) + T_{29}^0 e^{-\varepsilon_5 t}$ If we take t such that $e^{-\varepsilon_5 t} = \frac{1}{2}$ it results<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{2}\right)$, $t = log\frac{2}{\varepsilon_5}$ By taking now $\varepsilon_5$ sufficiently small one sees that $T_{29}$ is unbounded.<br>The same property holds for $T_{30}$ if $\lim_{t\to\infty}(b_{30}'')^{(5)}((G_{31})(t),t) = (b_{30}')^{(5)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations<br><br>Analogous inequalities hold also for $G_{33}$ , $G_{34}, T_{32}, T_{33}, T_{34}$ | 243 |
| It is now sufficient to take $\frac{(a_i)^{(6)}}{(\widehat{M}_{32})^{(6)}}$ , $\frac{(b_i)^{(6)}}{(\widehat{M}_{32})^{(6)}} < 1$ and to choose<br><br>$(\widehat{P}_{32})^{(6)}$ and $(\widehat{Q}_{32})^{(6)}$ large to have | 244 |
| $\frac{(a_i)^{(6)}}{(\widehat{M}_{32})^{(6)}}\left[(\widehat{P}_{32})^{(6)} + ((\widehat{P}_{32})^{(6)} + G_j^0)e^{-\left(\frac{(\widehat{P}_{32})^{(6)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{32})^{(6)}$ | 245 |

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| $\frac{(b_i)^{(6)}}{(\widehat{M}_{32})^{(6)}}\left[\left((\hat{Q}_{32})^{(6)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{32})^{(6)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{32})^{(6)}\right]\leq(\hat{Q}_{32})^{(6)}$ | 246 |
| In order that the operator $\mathcal{A}^{(6)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(6)}$ is a contraction with respect to the metric<br>$d\left(((G_{35})^{(1)},(T_{35})^{(1)}),((G_{35})^{(2)},(T_{35})^{(2)})\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{32})^{(6)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{32})^{(6)}t}\}$<br>Indeed if we denote<br>**Definition of** $(\widetilde{G_{35}}),(\widetilde{T_{35}})$ : $\left((\widetilde{G_{35}}),(\widetilde{T_{35}})\right)=\mathcal{A}^{(6)}\left((G_{35}),(T_{35})\right)$<br>It results<br>$\left\|\tilde{G}_{32}^{(1)}-\tilde{G}_i^{(2)}\right\|\leq\int_0^t(a_{32})^{(6)}\left\|G_{33}^{(1)}-G_{33}^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{(\widehat{M}_{32})^{(6)}s_{(32)}}ds_{(32)}+$<br>$\int_0^t\{(a'_{32})^{(6)}\left\|G_{32}^{(1)}-G_{32}^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}+$<br>$(a''_{32})^{(6)}\left(T_{33}^{(1)},s_{(32)}\right)\left\|G_{32}^{(1)}-G_{32}^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{(\widehat{M}_{32})^{(6)}s_{(32)}}+$<br>$G_{32}^{(2)}\left\|(a''_{32})^{(6)}\left(T_{33}^{(1)},s_{(32)}\right)-(a''_{32})^{(6)}\left(T_{33}^{(2)},s_{(32)}\right)\right\|\ e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{(\widehat{M}_{32})^{(6)}s_{(32)}}\}ds_{(32)}$<br>Where $s_{(32)}$ represents integrand that is integrated over the interval $[0,t]$<br>From the hypotheses it follows | 247 |
| $\left\|(G_{35})^{(1)}-(G_{35})^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}t}\leq$<br>$\frac{1}{(\widehat{M}_{32})^{(6)}}\left((a_{32})^{(6)}+(a'_{32})^{(6)}+(\widehat{A}_{32})^{(6)}+(\widehat{P}_{32})^{(6)}(\widehat{k}_{32})^{(6)}\right)d\left(((G_{35})^{(1)},(T_{35})^{(1)};\ (G_{35})^{(2)},(T_{35})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 248 |
| **Remark 26:** The fact that we supposed $(a''_{32})^{(6)}$ and $(b''_{32})^{(6)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{32})^{(6)}e^{(\widehat{M}_{32})^{(6)}t}$ *and* $(\widehat{Q}_{32})^{(6)}e^{(\widehat{M}_{32})^{(6)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a''_i)^{(6)}$ and $(b''_i)^{(6)}, i=32,33,34$ depend only on $T_{33}$ and respectively on $(G_{35})$(*and not on* $t$) and hypothesis can replaced by a usual Lipschitz condition. | 249 |

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| **Remark 27:** There does not exist any $t$ where $G_i(t) = 0$ *and* $T_i(t) = 0$<br><br>it results<br><br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i')^{(6)} - (a_i'')^{(6)}(T_{33}(s_{(32)}), s_{(32)})\} ds_{(32)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{(-(b_i')^{(6)} t)} > 0$ for t > 0 | 250 |
| **Definition of** $((\widehat{M}_{32})^{(6)})_1, ((\widehat{M}_{32})^{(6)})_2$ *and* $((\widehat{M}_{32})^{(6)})_3$ :<br><br>**Remark 28:** if $G_{32}$ is bounded, the same property have also $G_{33}$ *and* $G_{34}$ . indeed if<br><br>$G_{32} < (\widehat{M}_{32})^{(6)}$ it follows $\frac{dG_{33}}{dt} \leq ((\widehat{M}_{32})^{(6)})_1 - (a_{33}')^{(6)} G_{33}$ and by integrating<br><br>$G_{33} \leq ((\widehat{M}_{32})^{(6)})_2 = G_{33}^0 + 2(a_{33})^{(6)} ((\widehat{M}_{32})^{(6)})_1 / (a_{33}')^{(6)}$<br><br>In the same way , one can obtain<br><br>$G_{34} \leq ((\widehat{M}_{32})^{(6)})_3 = G_{34}^0 + 2(a_{34})^{(6)} ((\widehat{M}_{32})^{(6)})_2 / (a_{34}')^{(6)}$<br><br>If $G_{33}$ *or* $G_{34}$ is bounded, the same property follows for $G_{32}$ , $G_{34}$ and $G_{32}$ , $G_{33}$ respectively. | 251 |
| **Remark 29:** If $G_{32}$ *is* bounded, from below, the same property holds for $G_{33}$ *and* $G_{34}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{33}$ is bounded from below. | 252 |
| **Remark 30:** If $T_{32}$ is bounded from below and $\lim_{t\to\infty} ((b_i'')^{(6)}((G_{35})(t), t)) = (b_{33}')^{(6)}$ then $T_{33} \to \infty$.<br><br>**Definition of** $(m)^{(6)}$ and $\varepsilon_6$ :<br><br>Indeed let $t_6$ be so that for $t > t_6$<br><br>$(b_{33})^{(6)} - (b_i'')^{(6)}((G_{35})(t), t) < \varepsilon_6, T_{32}(t) > (m)^{(6)}$ | 253 |
| Then $\frac{dT_{33}}{dt} \geq (a_{33})^{(6)}(m)^{(6)} - \varepsilon_6 T_{33}$ which leads to<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{\varepsilon_6}\right)(1 - e^{-\varepsilon_6 t}) + T_{33}^0 e^{-\varepsilon_6 t}$ If we take t such that $e^{-\varepsilon_6 t} = \frac{1}{2}$ it results<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{2}\right)$, $t = log\frac{2}{\varepsilon_6}$ By taking now $\varepsilon_6$ sufficiently small one sees that $T_{33}$ is unbounded.<br>The same property holds for $T_{34}$ if $\lim_{t\to\infty} (b_{34}'')^{(6)}((G_{35})(t), t(t), t) = (b_{34}')^{(6)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 254 |
| Analogous inequalities hold also for $G_{37}$ , $G_{38}, T_{36}, T_{37}, T_{38}$<br><br>It is now sufficient to take $\frac{(a_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}$ , $\frac{(b_i)^{(7)}}{(\widehat{M}_{36})^{(7)}} < 1$ and to choose<br>$(\widehat{P}_{36})^{(7)}$ and $(\widehat{Q}_{36})^{(7)}$ large to have | 255 |

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| $\frac{(a_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}\left[(\widehat{P}_{36})^{(7)}+\left((\hat{P}_{36})^{(7)}+G_j^0\right)e^{-\left(\frac{(\hat{P}_{36})^{(7)}+G_j^0}{G_j^0}\right)}\right]\leq(\hat{P}_{36})^{(7)}$ | 256 |
| $\frac{(b_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}\left[\left((\hat{Q}_{36})^{(7)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{36})^{(7)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{36})^{(7)}\right]\leq(\hat{Q}_{36})^{(7)}$ | 257 |
| In order that the operator $\mathcal{A}^{(7)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(7)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{39})^{(1)},(T_{39})^{(1)}),((G_{39})^{(2)},(T_{39})^{(2)})\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{36})^{(7)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{36})^{(7)}t}\}$<br><br>Indeed if we denote<br><br>**Definition of** $\widetilde{(G_{39})},\widetilde{(T_{39})}:\left(\widetilde{(G_{39})},\widetilde{(T_{39})}\right)=\mathcal{A}^{(7)}((G_{39}),(T_{39}))$<br>It results<br><br>$\left\|\tilde{G}_{36}^{(1)}-\tilde{G}_i^{(2)}\right\|\leq\int_0^t(a_{36})^{(7)}\left\|G_{37}^{(1)}-G_{37}^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{(\widehat{M}_{36})^{(7)}s_{(36)}}\,ds_{(36)}+$<br>$\int_0^t\{(a_{36}')^{(7)}\left\|G_{36}^{(1)}-G_{36}^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}+$<br>$(a_{36}'')^{(7)}\left(T_{37}^{(1)},s_{(36)}\right)\left\|G_{36}^{(1)}-G_{36}^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{(\widehat{M}_{36})^{(7)}s_{(36)}}+$<br>$G_{36}^{(2)}\left\|(a_{36}'')^{(7)}\left(T_{37}^{(1)},s_{(36)}\right)-(a_{36}'')^{(7)}\left(T_{37}^{(2)},s_{(36)}\right)\right\|\ e^{-(\widehat{M}_{36})^{(7)}s_{(36)}}e^{(\widehat{M}_{36})^{(7)}s_{(36)}}\}ds_{(36)}$<br><br>Where $s_{(36)}$ represents integrand that is integrated over the interval $[0,t]$<br><br>From the hypotheses on it follows | 258 |
| $\left\|(G_{39})^{(1)}-(G_{39})^{(2)}\right\|e^{-(\widehat{M}_{36})^{(7)}t}\leq$<br>$\frac{1}{(\widehat{M}_{36})^{(7)}}\left((a_{36})^{(7)}+(a_{36}')^{(7)}+(\widehat{A}_{36})^{(7)}+(\widehat{P}_{36})^{(7)}(\widehat{k}_{36})^{(7)}\right)d\left(((G_{39})^{(1)},(T_{39})^{(1)};\ (G_{39})^{(2)},(T_{39})^{(2)})\right)$<br><br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 259 |
| **Remark 31:** The fact that we supposed $(a_{36}'')^{(7)}$ and $(b_{36}'')^{(7)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{36})^{(7)}e^{(\widehat{M}_{36})^{(7)}t}$ *and* $(\widehat{Q}_{36})^{(7)}e^{(\widehat{M}_{36})^{(7)}t}$ respectively of $\mathbb{R}_+$. | 260 |

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| If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(7)}$ and $(b_i^{''})^{(7)}, i = 36,37,38$ depend only on $T_{37}$ and respectively on $(G_{39})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 32:** There does not exist any $t$ where $G_i\,(t) = 0\ and\ T_i\,(t) = 0$<br><br>it results<br><br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(7)} - (a_i^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\} ds_{(36)}\right]} \geq 0$<br><br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(7)} t\right)} > 0$ for t > 0 | 261 |
| **Definition of** $\left((\widehat{M}_{36})^{(7)}\right)_1, \left((\widehat{M}_{36})^{(7)}\right)_2\ and\ \left((\widehat{M}_{36})^{(7)}\right)_3$ :<br><br>**Remark 33:** if $G_{36}$ is bounded, the same property have also $G_{37}\ and\ G_{38}$ . indeed if<br><br>$G_{36} < (\widehat{M}_{36})^{(7)}$ it follows $\frac{dG_{37}}{dt} \leq \left((\widehat{M}_{36})^{(7)}\right)_1 - (a_{37}^{'})^{(7)} G_{37}$ and by integrating<br><br>$G_{37} \leq \left((\widehat{M}_{36})^{(7)}\right)_2 = G_{37}^0 + 2(a_{37})^{(7)}\left((\widehat{M}_{36})^{(7)}\right)_1 / (a_{37}^{'})^{(7)}$<br><br>In the same way , one can obtain<br><br>$G_{38} \leq \left((\widehat{M}_{36})^{(7)}\right)_3 = G_{38}^0 + 2(a_{38})^{(7)}\left((\widehat{M}_{36})^{(7)}\right)_2 / (a_{38}^{'})^{(7)}$<br><br>If $G_{37}\ or\ G_{38}$ is bounded, the same property follows for $G_{36}$ , $G_{38}$ and $G_{36}$ , $G_{37}$ respectively. | 262 |
| **Remark 34:** If $G_{36}$ *is* bounded, from below, the same property holds for $G_{37}\ and\ G_{38}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{37}$ is bounded from below. | 263 |
| **Remark 35:** If $T_{36}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(7)}((G_{39})(t), t)) = (b_{37}^{'})^{(7)}$ then $T_{37} \to \infty$.<br><br>**Definition of** $(m)^{(7)}$ and $\varepsilon_7$ :<br><br>Indeed let $t_7$ be so that for $t > t_7$<br><br>$(b_{37})^{(7)} - (b_i^{''})^{(7)}((G_{39})(t), t) < \varepsilon_7, T_{36}\,(t) > (m)^{(7)}$ | 264 |
| Then $\frac{dT_{37}}{dt} \geq (a_{37})^{(7)}(m)^{(7)} - \varepsilon_7 T_{37}$ which leads to<br><br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{\varepsilon_7}\right)(1 - e^{-\varepsilon_7 t}) + T_{37}^0 e^{-\varepsilon_7 t}$ If we take t such that $e^{-\varepsilon_7 t} = \frac{1}{2}$ it results<br><br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{2}\right),\quad t = log\frac{2}{\varepsilon_7}$ By taking now $\varepsilon_7$ sufficiently small one sees that $T_{37}$ is unbounded. | 265 |

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| The same property holds for $T_{38}$ if $\lim_{t\to\infty}(b_{38}^{''})^{(7)}\left((G_{39})(t),t\right)=(b_{38}^{'})^{(7)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |
| It is now sufficient to take $\frac{(a_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}$ , $\frac{(b_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}<1$ and to choose<br>$(\widehat{P}_{40})^{(8)}$ and $(\widehat{Q}_{40})^{(8)}$ large to have | 266 |
| $\frac{(a_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}\left[(\widehat{P}_{40})^{(8)}+\left((\hat{P}_{40})^{(8)}+G_j^0\right)e^{-\left(\frac{(\hat{P}_{40})^{(8)}+G_j^0}{G_j^0}\right)}\right]\leq(\hat{P}_{40})^{(8)}$ | 267 |
| $\frac{(b_i)^{(8)}}{(\widehat{M}_{40})^{(8)}}\left[\left((\hat{Q}_{40})^{(8)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{40})^{(8)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{40})^{(8)}\right]\leq(\hat{Q}_{40})^{(8)}$ | 268 |
| In order that the operator $\mathcal{A}^{(8)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(8)}$ is a contraction with respect to the metric | |
| $d\left(((G_{43})^{(1)},(T_{43})^{(1)}),((G_{43})^{(2)},(T_{43})^{(2)})\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{40})^{(8)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{40})^{(8)}t}\}$ | 269 |
| Indeed if we denote<br><br>**Definition of** $\widetilde{(G_{43})},\widetilde{(T_{43})}$ : $\left(\widetilde{(G_{43})},\widetilde{(T_{43})}\right)=\mathcal{A}^{(8)}((G_{43}),(T_{43}))$ | 270 |
| It results<br><br>$\left\|\tilde{G}_{40}^{(1)}-\tilde{G}_i^{(2)}\right\|\leq\int_0^t(a_{40})^{(8)}\left\|G_{41}^{(1)}-G_{41}^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{(\widehat{M}_{40})^{(8)}s_{(40)}}ds_{(40)}+$<br>$\int_0^t\{(a_{40}^{'})^{(8)}\left\|G_{40}^{(1)}-G_{40}^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}+$<br>$(a_{40}^{''})^{(8)}\left(T_{41}^{(1)},s_{(40)}\right)\left\|G_{40}^{(1)}-G_{40}^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{(\widehat{M}_{40})^{(8)}s_{(40)}}+$<br>$G_{40}^{(2)}\left\\|(a_{40}^{''})^{(8)}\left(T_{41}^{(1)},s_{(40)}\right)-(a_{40}^{''})^{(8)}\left(T_{41}^{(2)},s_{(40)}\right)\right\\|\ e^{-(\widehat{M}_{40})^{(8)}s_{(40)}}e^{(\widehat{M}_{40})^{(8)}s_{(40)}}\}ds_{(40)}$ | 271 |
| Where $s_{(40)}$ represents integrand that is integrated over the interval $[0,t]$<br><br>From the hypotheses it follows | 272 |
| $\left\|(G_{43})^{(1)}-(G_{43})^{(2)}\right\|e^{-(\widehat{M}_{40})^{(8)}t}\leq$<br>$\frac{1}{(\widehat{M}_{40})^{(8)}}\left((a_{40})^{(8)}+(a_{40}^{'})^{(8)}+(\widehat{A}_{40})^{(8)}+(\widehat{P}_{40})^{(8)}(\widehat{k}_{40})^{(8)}\right)d\left(((G_{43})^{(1)},(T_{43})^{(1)};\ (G_{43})^{(2)},(T_{43})^{(2)})\right)$<br><br>And analogous inequalities for $G_i$ $and$ $T_i$. Taking into account the hypothesis the result follows | 273 |

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| **Remark 36:** The fact that we supposed $(a_{40}^{''})^{(8)}$ and $(b_{40}^{''})^{(8)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{40})^{(8)}e^{(\widehat{M}_{40})^{(8)}t}$ $and$ $(\widehat{Q}_{40})^{(8)}e^{(\widehat{M}_{40})^{(8)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(8)}$ and $(b_i^{''})^{(8)}, i = 40,41,42$ depend only on $T_{41}$ and respectively on $(G_{43})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 274 |
| **Remark 37** There does not exist any $t$ where $G_i\,(t) = 0\ and\ T_i\,(t) = 0$<br><br>it results<br><br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(8)} - (a_i^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\} ds_{(40)}\right]} \geq 0$<br><br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(8)}t\right)} > 0$ for t > 0 | 275 |
| **Definition of** $\left((\widehat{M}_{40})^{(8)}\right)_1, \left((\widehat{M}_{40})^{(8)}\right)_2\ and\ \left((\widehat{M}_{40})^{(8)}\right)_3$ :<br><br>**Remark 38:** if $G_{40}$ is bounded, the same property have also $G_{41}\ and\ G_{42}$ . indeed if<br><br>$G_{40} < (\widehat{M}_{40})^{(8)}$ it follows $\frac{dG_{41}}{dt} \leq \left((\widehat{M}_{40})^{(8)}\right)_1 - (a_{41}^{'})^{(8)} G_{41}$ and by integrating<br><br>$G_{41} \leq \left((\widehat{M}_{40})^{(8)}\right)_2 = G_{41}^0 + 2(a_{41})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_1 / (a_{41}^{'})^{(8)}$<br><br>In the same way , one can obtain<br><br>$G_{42} \leq \left((\widehat{M}_{40})^{(8)}\right)_3 = G_{42}^0 + 2(a_{42})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_2 / (a_{42}^{'})^{(8)}$<br><br>If $G_{41}\ or\ G_{42}$ is bounded, the same property follows for $G_{40}$ , $G_{42}$ and $G_{40}$ , $G_{41}$ respectively. | 276 |
| **Remark 39:** If $G_{40}$ $is$ bounded, from below, the same property holds for $G_{41}\ and\ G_{42}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{41}$ is bounded from below. | 277 |
| **Remark 40:** If $T_{40}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(8)}((G_{43})(t), t)) = (b_{41}^{'})^{(8)}$ then $T_{41} \to \infty$.<br><br>**Definition of** $(m)^{(8)}$ and $\varepsilon_8$ :<br><br>Indeed let $t_8$ be so that for $t > t_8$<br><br>$(b_{41})^{(8)} - (b_i^{''})^{(8)}\left((G_{43})(t), t\right) < \varepsilon_8, T_{40}\,(t) > (m)^{(8)}$ | 278 |

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| Then $\frac{dT_{41}}{dt} \geq (a_{41})^{(8)}(m)^{(8)} - \varepsilon_8 T_{41}$ which leads to<br><br>$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{\varepsilon_8}\right)(1 - e^{-\varepsilon_8 t}) + T_{41}^0 e^{-\varepsilon_8 t}$ If we take t such that $e^{-\varepsilon_8 t} = \frac{1}{2}$ it results<br><br>$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{2}\right), \quad t = log\frac{2}{\varepsilon_8}$ By taking now $\varepsilon_8$ sufficiently small one sees that $T_{41}$ is unbounded.<br>The same property holds for $T_{42}$ if $\lim_{t\to\infty}(b_{42}^{''})^{(8)}\left((G_{43})(t), t(t), t\right) = (b_{42}^{'})^{(8)}$ | 279 |
| It is now sufficient to take $\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}, \frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}} < 1$ and to choose $(\widehat{P}_{44})^{(9)}$ and $(\widehat{Q}_{44})^{(9)}$ large to have | 279 A |
| $\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[(\widehat{P}_{44})^{(9)} + \left((\widehat{P}_{44})^{(9)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{44})^{(9)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{44})^{(9)}$ | |
| $\frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[\left((\widehat{Q}_{44})^{(9)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{44})^{(9)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{44})^{(9)}\right] \leq (\widehat{Q}_{44})^{(9)}$ | |
| In order that the operator $\mathcal{A}^{(9)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying 39,35,36 into itself | |
| The operator $\mathcal{A}^{(9)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{47})^{(1)}, (T_{47})^{(1)}), ((G_{47})^{(2)}, (T_{47})^{(2)})\right) =$<br>$sup_i\{max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t) - G_i^{(2)}(t)\right|e^{-(\widehat{M}_{44})^{(9)}t}, max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t) - T_i^{(2)}(t)\right|e^{-(\widehat{M}_{44})^{(9)}t}\}$<br><br>Indeed if we denote<br><br>**Definition of** $\widetilde{(G_{47})}, \widetilde{(T_{47})} : \left(\widetilde{(G_{47})}, \widetilde{(T_{47})}\right) = \mathcal{A}^{(9)}\left((G_{47}), (T_{47})\right)$<br>It results<br><br>$\left\|\tilde{G}_{44}^{(1)} - \tilde{G}_i^{(2)}\right| \leq \int_0^t (a_{44})^{(9)}\left|G_{45}^{(1)} - G_{45}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\,ds_{(44)} +$<br>$\int_0^t\{(a_{44}^{'})^{(9)}\left|G_{44}^{(1)} - G_{44}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{-(\widehat{M}_{44})^{(9)}s_{(44)}} +$<br>$(a_{44}^{''})^{(9)}\left(T_{45}^{(1)}, s_{(44)}\right)\left|G_{44}^{(1)} - G_{44}^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}} +$<br>$G_{44}^{(2)}\left|(a_{44}^{''})^{(9)}\left(T_{45}^{(1)}, s_{(44)}\right) - (a_{44}^{''})^{(9)}\left(T_{45}^{(2)}, s_{(44)}\right)\right| \; e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\}ds_{(44)}$<br><br>Where $s_{(44)}$ represents integrand that is integrated over the interval $[0, t]$<br><br>From the hypotheses on 45,46,47,28 and 29 it follows | |
| $\left|(G_{47})^{(1)} - G^{(2)}\right|e^{-(\widehat{M}_{44})^{(9)}t} \leq$<br>$\frac{1}{(\widehat{M}_{44})^{(9)}}\left((a_{44})^{(9)} + (a_{44}^{'})^{(9)} + (\widehat{A}_{44})^{(9)} + (\widehat{P}_{44})^{(9)}(\widehat{k}_{44})^{(9)}\right)d\left(((G_{47})^{(1)}, (T_{47})^{(1)}; (G_{47})^{(2)}, (T_{47})^{(2)})\right)$ | |

And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis (39,35,36) the result follows

**Remark 41:** The fact that we supposed $(a_{44}^{''})^{(9)}$ and $(b_{44}^{''})^{(9)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{44})^{(9)}e^{(\widehat{M}_{44})^{(9)}t}$ *and* $(\widehat{Q}_{44})^{(9)}e^{(\widehat{M}_{44})^{(9)}t}$ respectively of $\mathbb{R}_+$.

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(9)}$ and $(b_i^{''})^{(9)}, i = 44,45,46$ depend only on $T_{45}$ and respectively on $(G_{47})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition.

**Remark 42:** There does not exist any $t$ where $G_i\,(t) = 0\ and\ T_i\,(t) = 0$

From 99 to 44 it results

$$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(9)} - (a_i^{''})^{(9)}(T_{45}(s_{(44)}), s_{(44)})\} ds_{(44)}\right]} \geq 0$$

$$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(9)}t\right)} > 0 \quad \text{for t} > 0$$

**Definition of** $\left((\widehat{M}_{44})^{(9)}\right)_1, \left((\widehat{M}_{44})^{(9)}\right)_2\ and\ \left((\widehat{M}_{44})^{(9)}\right)_3$ :

**Remark 43:** if $G_{44}$ is bounded, the same property have also $G_{45}\ and\ G_{46}$ . indeed if

$G_{44} < (\widehat{M}_{44})^{(9)}$ it follows $\frac{dG_{45}}{dt} \leq \left((\widehat{M}_{44})^{(9)}\right)_1 - (a_{45}^{'})^{(9)}G_{45}$ and by integrating

$$G_{45} \leq \left((\widehat{M}_{44})^{(9)}\right)_2 = G_{45}^0 + 2(a_{45})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_1/(a_{45}^{'})^{(9)}$$

In the same way , one can obtain

$$G_{46} \leq \left((\widehat{M}_{44})^{(9)}\right)_3 = G_{46}^0 + 2(a_{46})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_2/(a_{46}^{'})^{(9)}$$

If $G_{45}\ or\ G_{46}$ is bounded, the same property follows for $G_{44}$ , $G_{46}$ and $G_{44}$ , $G_{45}$ respectively.

**Remark 44:** If $G_{44}$ *is* bounded, from below, the same property holds for $G_{45}\ and\ G_{46}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{45}$ is bounded from below.

**Remark 45:** If $T_{44}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(9)}\left((G_{47})(t), t\right) = (b_{45}^{'})^{(9)}$ then $T_{45} \to \infty$.

**Definition of** $(m)^{(9)}$ and $\varepsilon_9$ :

Indeed let $t_9$ be so that for $t > t_9$

$$(b_{45})^{(9)} - (b_i^{''})^{(9)}\left((G_{47})(t), t\right) < \varepsilon_9, T_{44}\,(t) > (m)^{(9)}$$

Then $\frac{dT_{45}}{dt} \geq (a_{45})^{(9)}(m)^{(9)} - \varepsilon_9 T_{45}$ which leads to

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{\varepsilon_9}\right)(1 - e^{-\varepsilon_9 t}) + T_{45}^0 e^{-\varepsilon_9 t}$ If we take t such that $e^{-\varepsilon_9 t} = \frac{1}{2}$ it results

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{2}\right)$, $t = log\frac{2}{\varepsilon_9}$ By taking now $\varepsilon_9$ sufficiently small one sees that $T_{45}$ is unbounded. The same property holds for $T_{46}$ if $\lim_{t\to\infty}(b''_{46})^{(9)}\left((G_{47})(t),t\right) = (b'_{46})^{(9)}$
We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 92

**Behavior of the solutions of equation** 280

**Theorem** If we denote and define

**Definition of** $(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ :

$(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ four constants satisfying
$-(\sigma_2)^{(1)} \leq -(a'_{13})^{(1)} + (a'_{14})^{(1)} - (a''_{13})^{(1)}(T_{14}, t) + (a''_{14})^{(1)}(T_{14}, t) \leq -(\sigma_1)^{(1)}$

$-(\tau_2)^{(1)} \leq -(b'_{13})^{(1)} + (b'_{14})^{(1)} - (b''_{13})^{(1)}(G,t) - (b''_{14})^{(1)}(G,t) \leq -(\tau_1)^{(1)}$

**Definition of** $(\nu_1)^{(1)}, (\nu_2)^{(1)}, (u_1)^{(1)}, (u_2)^{(1)}, \nu^{(1)}, u^{(1)}$ : 281

By $(\nu_1)^{(1)} > 0$, $(\nu_2)^{(1)} < 0$ and respectively $(u_1)^{(1)} > 0$, $(u_2)^{(1)} < 0$ the roots of the equations $(a_{14})^{(1)}\left(\nu^{(1)}\right)^2 + (\sigma_1)^{(1)}\nu^{(1)} - (a_{13})^{(1)} = 0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2 + (\tau_1)^{(1)}u^{(1)} - (b_{13})^{(1)} = 0$

**Definition of** $(\bar{\nu}_1)^{(1)},, (\bar{\nu}_2)^{(1)}, (\bar{u}_1)^{(1)}, (\bar{u}_2)^{(1)}$ : 282

By $(\bar{\nu}_1)^{(1)} > 0$, $(\bar{\nu}_2)^{(1)} < 0$ and respectively $(\bar{u}_1)^{(1)} > 0$, $(\bar{u}_2)^{(1)} < 0$ the roots of the equations $(a_{14})^{(1)}\left(\nu^{(1)}\right)^2 + (\sigma_2)^{(1)}\nu^{(1)} - (a_{13})^{(1)} = 0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2 + (\tau_2)^{(1)}u^{(1)} - (b_{13})^{(1)} = 0$

**Definition of** $(m_1)^{(1)}, (m_2)^{(1)}, (\mu_1)^{(1)}, (\mu_2)^{(1)}, (\nu_0)^{(1)}$ :- 283

If we define $(m_1)^{(1)}, (m_2)^{(1)}, (\mu_1)^{(1)}, (\mu_2)^{(1)}$ by
$(m_2)^{(1)} = (\nu_0)^{(1)}, (m_1)^{(1)} = (\nu_1)^{(1)}$, $if\ (\nu_0)^{(1)} < (\nu_1)^{(1)}$

$(m_2)^{(1)} = (\nu_1)^{(1)}, (m_1)^{(1)} = (\bar{\nu}_1)^{(1)}, if\ (\nu_1)^{(1)} < (\nu_0)^{(1)} < (\bar{\nu}_1)^{(1)}$,

and $\boxed{(\nu_0)^{(1)} = \frac{G^0_{13}}{G^0_{14}}}$

$(m_2)^{(1)} = (\nu_1)^{(1)}, (m_1)^{(1)} = (\nu_0)^{(1)}$, $if\ (\bar{\nu}_1)^{(1)} < (\nu_0)^{(1)}$

and analogously 284

$(\mu_2)^{(1)} = (u_0)^{(1)}, (\mu_1)^{(1)} = (u_1)^{(1)}$, $if\ (u_0)^{(1)} < (u_1)^{(1)}$

$(\mu_2)^{(1)} = (u_1)^{(1)}, (\mu_1)^{(1)} = (\bar{u}_1)^{(1)}, if\ (u_1)^{(1)} < (u_0)^{(1)} < (\bar{u}_1)^{(1)}$,

and $\boxed{(u_0)^{(1)} = \frac{T^0_{13}}{T^0_{14}}}$

$(\mu_2)^{(1)} = (u_1)^{(1)}, (\mu_1)^{(1)} = (u_0)^{(1)}, if\ (\bar{u}_1)^{(1)} < (u_0)^{(1)}$ where $(u_1)^{(1)}, (\bar{u}_1)^{(1)}$

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| are defined | |
| Then the solution of global equations satisfies the inequalities<br>$G_{13}^0 e^{((S_1)^{(1)}-(p_{13})^{(1)})t} \le G_{13}(t) \le G_{13}^0 e^{(S_1)^{(1)}t}$<br>where $(p_i)^{(1)}$ is defined by equation<br>$\frac{1}{(m_1)^{(1)}} G_{13}^0 e^{((S_1)^{(1)}-(p_{13})^{(1)})t} \le G_{14}(t) \le \frac{1}{(m_2)^{(1)}} G_{13}^0 e^{(S_1)^{(1)}t}$ | 285 |
| $(\frac{(a_{15})^{(1)}G_{13}^0}{(m_1)^{(1)}((S_1)^{(1)}-(p_{13})^{(1)}-(S_2)^{(1)})}\left[e^{((S_1)^{(1)}-(p_{13})^{(1)})t} - e^{-(S_2)^{(1)}t}\right] + G_{15}^0 e^{-(S_2)^{(1)}t} \le G_{15}(t) \le$<br>$\frac{(a_{15})^{(1)}G_{13}^0}{(m_2)^{(1)}((S_1)^{(1)}-(a'_{15})^{(1)})}[e^{(S_1)^{(1)}t} - e^{-(a'_{15})^{(1)}t}] + G_{15}^0 e^{-(a'_{15})^{(1)}t})$ | 286 |
| $\boxed{T_{13}^0 e^{(R_1)^{(1)}t} \le T_{13}(t) \le T_{13}^0 e^{((R_1)^{(1)}+(r_{13})^{(1)})t}}$ | 287 |
| $\frac{1}{(\mu_1)^{(1)}} T_{13}^0 e^{(R_1)^{(1)}t} \le T_{13}(t) \le \frac{1}{(\mu_2)^{(1)}} T_{13}^0 e^{((R_1)^{(1)}+(r_{13})^{(1)})t}$ | 288 |
| $\frac{(b_{15})^{(1)}T_{13}^0}{(\mu_1)^{(1)}((R_1)^{(1)}-(b'_{15})^{(1)})}\left[e^{(R_1)^{(1)}t} - e^{-(b'_{15})^{(1)}t}\right] + T_{15}^0 e^{-(b'_{15})^{(1)}t} \le T_{15}(t) \le$<br>$\frac{(a_{15})^{(1)}T_{13}^0}{(\mu_2)^{(1)}((R_1)^{(1)}+(r_{13})^{(1)}+(R_2)^{(1)})}\left[e^{((R_1)^{(1)}+(r_{13})^{(1)})t} - e^{-(R_2)^{(1)}t}\right] + T_{15}^0 e^{-(R_2)^{(1)}t}$ | 289 |
| **Definition of** $(S_1)^{(1)}, (S_2)^{(1)}, (R_1)^{(1)}, (R_2)^{(1)}$:-<br>Where $(S_1)^{(1)} = (a_{13})^{(1)}(m_2)^{(1)} - (a'_{13})^{(1)}$<br>$(S_2)^{(1)} = (a_{15})^{(1)} - (p_{15})^{(1)}$<br>$(R_1)^{(1)} = (b_{13})^{(1)}(\mu_2)^{(1)} - (b'_{13})^{(1)}$<br>$(R_2)^{(1)} = (b'_{15})^{(1)} - (r_{15})^{(1)}$ | 290 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define | 291 |
| **Definition of** $(\sigma_1)^{(2)}, (\sigma_2)^{(2)}, (\tau_1)^{(2)}, (\tau_2)^{(2)}$ :<br>$(\sigma_1)^{(2)}, (\sigma_2)^{(2)}, (\tau_1)^{(2)}, (\tau_2)^{(2)}$ four constants satisfying | 292 |
| $-(\sigma_2)^{(2)} \le -(a'_{16})^{(2)} + (a'_{17})^{(2)} - (a''_{16})^{(2)}(T_{17}, t) + (a''_{17})^{(2)}(T_{17}, t) \le -(\sigma_1)^{(2)}$ | 293 |
| $-(\tau_2)^{(2)} \le -(b'_{16})^{(2)} + (b'_{17})^{(2)} - (b''_{16})^{(2)}((G_{19}), t) - (b''_{17})^{(2)}((G_{19}), t) \le -(\tau_1)^{(2)}$ | 294 |
| **Definition of** $(v_1)^{(2)}, (v_2)^{(2)}, (u_1)^{(2)}, (u_2)^{(2)}$ : | 295 |

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| By $(\nu_1)^{(2)} > 0$ , $(\nu_2)^{(2)} < 0$ and respectively $(u_1)^{(2)} > 0$ , $(u_2)^{(2)} < 0$ the roots | 296 |
| of the equations $(a_{17})^{(2)}\left(\nu^{(2)}\right)^2 + (\sigma_1)^{(2)}\nu^{(2)} - (a_{16})^{(2)} = 0$ | 297 |
| and $(b_{14})^{(2)}\left(u^{(2)}\right)^2 + (\tau_1)^{(2)}u^{(2)} - (b_{16})^{(2)} = 0$ and | 298 |
| **Definition of** $(\bar{\nu}_1)^{(2)}$,, $(\bar{\nu}_2)^{(2)}, (\bar{u}_1)^{(2)}, (\bar{u}_2)^{(2)}$ : | 299 |
| By $(\bar{\nu}_1)^{(2)} > 0$ , $(\bar{\nu}_2)^{(2)} < 0$ and respectively $(\bar{u}_1)^{(2)} > 0$ , $(\bar{u}_2)^{(2)} < 0$ the | 300 |
| roots of the equations $(a_{17})^{(2)}\left(\nu^{(2)}\right)^2 + (\sigma_2)^{(2)}\nu^{(2)} - (a_{16})^{(2)} = 0$ | 301 |
| and $(b_{17})^{(2)}\left(u^{(2)}\right)^2 + (\tau_2)^{(2)}u^{(2)} - (b_{16})^{(2)} = 0$ | 302 |
| **Definition of** $(m_1)^{(2)}$ , $(m_2)^{(2)}$ , $(\mu_1)^{(2)}, (\mu_2)^{(2)}$ :- | 303 |
| If we define $(m_1)^{(2)}$ , $(m_2)^{(2)}$ , $(\mu_1)^{(2)}, (\mu_2)^{(2)}$ by | 304 |
| $(m_2)^{(2)} = (\nu_0)^{(2)}, (m_1)^{(2)} = (\nu_1)^{(2)}$, ***if*** $(\nu_0)^{(2)} < (\nu_1)^{(2)}$ | 305 |
| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ , ***if*** $(\nu_1)^{(2)} < (\nu_0)^{(2)} < (\bar{\nu}_1)^{(2)}$, and $\boxed{(\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0}}$ | 306 |
| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\nu_0)^{(2)}$, ***if*** $(\bar{\nu}_1)^{(2)} < (\nu_0)^{(2)}$ | 307 |
| and analogously $(\mu_2)^{(2)} = (u_0)^{(2)}, (\mu_1)^{(2)} = (u_1)^{(2)}$, ***if*** $(u_0)^{(2)} < (u_1)^{(2)}$ $(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (\bar{u}_1)^{(2)}$ , ***if*** $(u_1)^{(2)} < (u_0)^{(2)} < (\bar{u}_1)^{(2)}$, and $\boxed{(u_0)^{(2)} = \frac{T_{16}^0}{T_{17}^0}}$ | 308 |
| $(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (u_0)^{(2)}$, ***if*** $(\bar{u}_1)^{(2)} < (u_0)^{(2)}$ | 309 |
| Then the solution of global equations satisfies the inequalities $G_{16}^0 e^{((S_1)^{(2)} - (p_{16})^{(2)})t} \le G_{16}(t) \le G_{16}^0 e^{(S_1)^{(2)}t}$ | 310 |
| $(p_i)^{(2)}$ is defined by equation | |
| $\frac{1}{(m_1)^{(2)}} G_{16}^0 e^{((S_1)^{(2)} - (p_{16})^{(2)})t} \le G_{17}(t) \le \frac{1}{(m_2)^{(2)}} G_{16}^0 e^{(S_1)^{(2)}t}$ | 311 |
| $\left(\frac{(a_{18})^{(2)} G_{16}^0}{(m_1)^{(2)}((S_1)^{(2)} - (p_{16})^{(2)} - (S_2)^{(2)})}\left[e^{((S_1)^{(2)} - (p_{16})^{(2)})t} - e^{-(S_2)^{(2)}t}\right] + G_{18}^0 e^{-(S_2)^{(2)}t} \le G_{18}(t) \le\right.$ | 312 |

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| $\frac{(a_{18})^{(2)}G_{16}^{0}}{(m_2)^{(2)}((S_1)^{(2)}-(a_{18}')^{(2)})}[e^{(S_1)^{(2)}t}-e^{-(a_{18}')^{(2)}t}]+G_{18}^{0}e^{-(a_{18}')^{(2)}t})$ | |
| $\boxed{T_{16}^{0}e^{(R_1)^{(2)}t}\le T_{16}(t)\le T_{16}^{0}e^{((R_1)^{(2)}+(r_{16})^{(2)})t}}$ | 313 |
| $\frac{1}{(\mu_1)^{(2)}}T_{16}^{0}e^{(R_1)^{(2)}t}\le T_{16}(t)\le\frac{1}{(\mu_2)^{(2)}}T_{16}^{0}e^{((R_1)^{(2)}+(r_{16})^{(2)})t}$ | 314 |
| $\frac{(b_{18})^{(2)}T_{16}^{0}}{(\mu_1)^{(2)}((R_1)^{(2)}-(b_{18}')^{(2)})}\left[e^{(R_1)^{(2)}t}-e^{-(b_{18}')^{(2)}t}\right]+T_{18}^{0}e^{-(b_{18}')^{(2)}t}\le T_{18}(t)\le$ $\frac{(a_{18})^{(2)}T_{16}^{0}}{(\mu_2)^{(2)}((R_1)^{(2)}+(r_{16})^{(2)}+(R_2)^{(2)})}\left[e^{((R_1)^{(2)}+(r_{16})^{(2)})t}-e^{-(R_2)^{(2)}t}\right]+T_{18}^{0}e^{-(R_2)^{(2)}t}$ | 315 |
| **Definition of** $(S_1)^{(2)},(S_2)^{(2)},(R_1)^{(2)},(R_2)^{(2)}$:- | 316 |
| Where $(S_1)^{(2)}=(a_{16})^{(2)}(m_2)^{(2)}-(a_{16}')^{(2)}$ $(S_2)^{(2)}=(a_{18})^{(2)}-(p_{18})^{(2)}$ | 317 |
| $(R_1)^{(2)}=(b_{16})^{(2)}(\mu_2)^{(1)}-(b_{16}')^{(2)}$ $(R_2)^{(2)}=(b_{18}')^{(2)}-(r_{18})^{(2)}$ | 318 |
| **Behavior of the solutions** **Theorem 3:** If we denote and define **Definition of** $(\sigma_1)^{(3)},(\sigma_2)^{(3)},(\tau_1)^{(3)},(\tau_2)^{(3)}$: $(\sigma_1)^{(3)},(\sigma_2)^{(3)},(\tau_1)^{(3)},(\tau_2)^{(3)}$ four constants satisfying $-(\sigma_2)^{(3)}\le-(a_{20}')^{(3)}+(a_{21}')^{(3)}-(a_{20}'')^{(3)}(T_{21},t)+(a_{21}'')^{(3)}(T_{21},t)\le-(\sigma_1)^{(3)}$ $-(\tau_2)^{(3)}\le-(b_{20}')^{(3)}+(b_{21}')^{(3)}-(b_{20}'')^{(3)}(G_{23},t)-(b_{21}'')^{(3)}((G_{23}),t)\le-(\tau_1)^{(3)}$ | 319 |
| **Definition of** $(v_1)^{(3)},(v_2)^{(3)},(u_1)^{(3)},(u_2)^{(3)}$: By $(v_1)^{(3)}>0,(v_2)^{(3)}<0$ and respectively $(u_1)^{(3)}>0,(u_2)^{(3)}<0$ the roots of the equations $(a_{21})^{(3)}(v^{(3)})^2+(\sigma_1)^{(3)}v^{(3)}-(a_{20})^{(3)}=0$ and $(b_{21})^{(3)}(u^{(3)})^2+(\tau_1)^{(3)}u^{(3)}-(b_{20})^{(3)}=0$ and By $(\bar{v}_1)^{(3)}>0,(\bar{v}_2)^{(3)}<0$ and respectively $(\bar{u}_1)^{(3)}>0,(\bar{u}_2)^{(3)}<0$ the roots of the equations $(a_{21})^{(3)}(v^{(3)})^2+(\sigma_2)^{(3)}v^{(3)}-(a_{20})^{(3)}=0$ and $(b_{21})^{(3)}(u^{(3)})^2+(\tau_2)^{(3)}u^{(3)}-(b_{20})^{(3)}=0$ | 320 |
| **Definition of** $(m_1)^{(3)},(m_2)^{(3)},(\mu_1)^{(3)},(\mu_2)^{(3)}$:- | 321 |

If we define $(m_1)^{(3)}, (m_2)^{(3)}, (\mu_1)^{(3)}, (\mu_2)^{(3)}$ by

$(m_2)^{(3)} = (\nu_0)^{(3)}, (m_1)^{(3)} = (\nu_1)^{(3)}$, ***if*** $(\nu_0)^{(3)} < (\nu_1)^{(3)}$

$(m_2)^{(3)} = (\nu_1)^{(3)}, (m_1)^{(3)} = (\bar{\nu}_1)^{(3)}$, ***if*** $(\nu_1)^{(3)} < (\nu_0)^{(3)} < (\bar{\nu}_1)^{(3)}$,

and $\boxed{(\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0}}$

$(m_2)^{(3)} = (\nu_1)^{(3)}, (m_1)^{(3)} = (\nu_0)^{(3)}$, ***if*** $(\bar{\nu}_1)^{(3)} < (\nu_0)^{(3)}$

and analogously

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$(\mu_2)^{(3)} = (u_0)^{(3)}, (\mu_1)^{(3)} = (u_1)^{(3)}$, ***if*** $(u_0)^{(3)} < (u_1)^{(3)}$

$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (\bar{u}_1)^{(3)}$, ***if*** $(u_1)^{(3)} < (u_0)^{(3)} < (\bar{u}_1)^{(3)}$, and $\boxed{(u_0)^{(3)} = \frac{T_{20}^0}{T_{21}^0}}$

$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (u_0)^{(3)}$, ***if*** $(\bar{u}_1)^{(3)} < (u_0)^{(3)}$

Then the solution of global equations satisfies the inequalities

$$G_{20}^0 e^{((S_1)^{(3)} - (p_{20})^{(3)})t} \le G_{20}(t) \le G_{20}^0 e^{(S_1)^{(3)} t}$$

$(p_i)^{(3)}$ is defined by equation

$$\frac{1}{(m_1)^{(3)}} G_{20}^0 e^{((S_1)^{(3)} - (p_{20})^{(3)})t} \le G_{21}(t) \le \frac{1}{(m_2)^{(3)}} G_{20}^0 e^{(S_1)^{(3)} t} \quad (323)$$

$$\left(\frac{(a_{22})^{(3)} G_{20}^0}{(m_1)^{(3)}((S_1)^{(3)} - (p_{20})^{(3)} - (S_2)^{(3)})}\left[e^{((S_1)^{(3)} - (p_{20})^{(3)})t} - e^{-(S_2)^{(3)} t}\right] + G_{22}^0 e^{-(S_2)^{(3)} t} \le G_{22}(t) \le \right.$$

$$\left. \frac{(a_{22})^{(3)} G_{20}^0}{(m_2)^{(3)}((S_1)^{(3)} - (a_{22}')^{(3)})}\left[e^{(S_1)^{(3)} t} - e^{-(a_{22}')^{(3)} t}\right] + G_{22}^0 e^{-(a_{22}')^{(3)} t}\right) \quad (324)$$

$$\boxed{T_{20}^0 e^{(R_1)^{(3)} t} \le T_{20}(t) \le T_{20}^0 e^{((R_1)^{(3)} + (r_{20})^{(3)})t}} \quad (325)$$

$$\frac{1}{(\mu_1)^{(3)}} T_{20}^0 e^{(R_1)^{(3)} t} \le T_{20}(t) \le \frac{1}{(\mu_2)^{(3)}} T_{20}^0 e^{((R_1)^{(3)} + (r_{20})^{(3)})t} \quad (326)$$

$$\frac{(b_{22})^{(3)} T_{20}^0}{(\mu_1)^{(3)}((R_1)^{(3)} - (b_{22}')^{(3)})}\left[e^{(R_1)^{(3)} t} - e^{-(b_{22}')^{(3)} t}\right] + T_{22}^0 e^{-(b_{22}')^{(3)} t} \le T_{22}(t) \le$$

$$\frac{(a_{22})^{(3)} T_{20}^0}{(\mu_2)^{(3)}((R_1)^{(3)} + (r_{20})^{(3)} + (R_2)^{(3)})}\left[e^{((R_1)^{(3)} + (r_{20})^{(3)})t} - e^{-(R_2)^{(3)} t}\right] + T_{22}^0 e^{-(R_2)^{(3)} t} \quad (327)$$

**Definition of** $(S_1)^{(3)}, (S_2)^{(3)}, (R_1)^{(3)}, (R_2)^{(3)}$:- (328)

Where $(S_1)^{(3)} = (a_{20})^{(3)}(m_2)^{(3)} - (a_{20}')^{(3)}$

$(S_2)^{(3)} = (a_{22})^{(3)} - (p_{22})^{(3)}$

$$(R_1)^{(3)} = (b_{20})^{(3)}(\mu_2)^{(3)} - (b_{20}^{'})^{(3)}$$

$$(R_2)^{(3)} = (b_{22}^{'})^{(3)} - (r_{22})^{(3)}$$

**Behavior of the solutions of equation**

**Theorem:** If we denote and define

**Definition of** $(\sigma_1)^{(4)}, (\sigma_2)^{(4)}, (\tau_1)^{(4)}, (\tau_2)^{(4)}$ :

$(\sigma_1)^{(4)}, (\sigma_2)^{(4)}, (\tau_1)^{(4)}, (\tau_2)^{(4)}$ four constants satisfying

$$-(\sigma_2)^{(4)} \le -(a_{24}^{'})^{(4)} + (a_{25}^{'})^{(4)} - (a_{24}^{''})^{(4)}(T_{25}, t) + (a_{25}^{''})^{(4)}(T_{25}, t) \le -(\sigma_1)^{(4)}$$

$$-(\tau_2)^{(4)} \le -(b_{24}^{'})^{(4)} + (b_{25}^{'})^{(4)} - (b_{24}^{''})^{(4)}\big((G_{27}), t\big) - (b_{25}^{''})^{(4)}\big((G_{27}), t\big) \le -(\tau_1)^{(4)}$$

**Definition of** $(v_1)^{(4)}, (v_2)^{(4)}, (u_1)^{(4)}, (u_2)^{(4)}, v^{(4)}, u^{(4)}$ : 329

By $(v_1)^{(4)} > 0$, $(v_2)^{(4)} < 0$ and respectively $(u_1)^{(4)} > 0$, $(u_2)^{(4)} < 0$ the roots of the equations $(a_{25})^{(4)}\big(v^{(4)}\big)^2 + (\sigma_1)^{(4)}v^{(4)} - (a_{24})^{(4)} = 0$ and $(b_{25})^{(4)}\big(u^{(4)}\big)^2 + (\tau_1)^{(4)}u^{(4)} - (b_{24})^{(4)} = 0$ and

**Definition of** $(\bar{v}_1)^{(4)}, , (\bar{v}_2)^{(4)}, (\bar{u}_1)^{(4)}, (\bar{u}_2)^{(4)}$ : 330

By $(\bar{v}_1)^{(4)} > 0$, $(\bar{v}_2)^{(4)} < 0$ and respectively $(\bar{u}_1)^{(4)} > 0$, $(\bar{u}_2)^{(4)} < 0$ the roots of the equations $(a_{25})^{(4)}\big(v^{(4)}\big)^2 + (\sigma_2)^{(4)}v^{(4)} - (a_{24})^{(4)} = 0$ and $(b_{25})^{(4)}\big(u^{(4)}\big)^2 + (\tau_2)^{(4)}u^{(4)} - (b_{24})^{(4)} = 0$

**Definition of** $(m_1)^{(4)}, (m_2)^{(4)}, (\mu_1)^{(4)}, (\mu_2)^{(4)}, (v_0)^{(4)}$ :-

If we define $(m_1)^{(4)}, (m_2)^{(4)}, (\mu_1)^{(4)}, (\mu_2)^{(4)}$ by

$$(m_2)^{(4)} = (v_0)^{(4)}, (m_1)^{(4)} = (v_1)^{(4)}, \ \boldsymbol{if}\ (v_0)^{(4)} < (v_1)^{(4)}$$

$$(m_2)^{(4)} = (v_1)^{(4)}, (m_1)^{(4)} = (\bar{v}_1)^{(4)}, \boldsymbol{if}\ (v_4)^{(4)} < (v_0)^{(4)} < (\bar{v}_1)^{(4)},$$

and $\boxed{(v_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0}}$

$$(m_2)^{(4)} = (v_4)^{(4)}, (m_1)^{(4)} = (v_0)^{(4)}, \ \boldsymbol{if}\ (\bar{v}_4)^{(4)} < (v_0)^{(4)}$$

and analogously 331

$$(\mu_2)^{(4)} = (u_0)^{(4)}, (\mu_1)^{(4)} = (u_1)^{(4)}, \ \boldsymbol{if}\ (u_0)^{(4)} < (u_1)^{(4)}$$

$$(\mu_2)^{(4)} = (u_1)^{(4)}, (\mu_1)^{(4)} = (\bar{u}_1)^{(4)}, \boldsymbol{if}\ (u_1)^{(4)} < (u_0)^{(4)} < (\bar{u}_1)^{(4)},$$

and $\boxed{(u_0)^{(4)} = \frac{T_{24}^0}{T_{25}^0}}$

$$(\mu_2)^{(4)} = (u_1)^{(4)}, (\mu_1)^{(4)} = (u_0)^{(4)}, \boldsymbol{if}\ (\bar{u}_1)^{(4)} < (u_0)^{(4)} \text{ where } (u_1)^{(4)}, (\bar{u}_1)^{(4)}$$

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| Then the solution of global equations satisfies the inequalities<br>$G_{24}^0 e^{((S_1)^{(4)}-(p_{24})^{(4)})t} \le G_{24}(t) \le G_{24}^0 e^{(S_1)^{(4)}t}$<br>where $(p_i)^{(4)}$ is defined by equation | 332 |
| $\frac{1}{(m_1)^{(4)}} G_{24}^0 e^{((S_1)^{(4)}-(p_{24})^{(4)})t} \le G_{25}(t) \le \frac{1}{(m_2)^{(4)}} G_{24}^0 e^{(S_1)^{(4)}t}$ | 333 |
| $\left(\frac{(a_{26})^{(4)} G_{24}^0}{(m_1)^{(4)}((S_1)^{(4)}-(p_{24})^{(4)}-(S_2)^{(4)})}\left[e^{((S_1)^{(4)}-(p_{24})^{(4)})t} - e^{-(S_2)^{(4)}t}\right] + G_{26}^0 e^{-(S_2)^{(4)}t} \le G_{26}(t) \le$<br>*(a26)4G240(m2)4(S1)4−(a26′)4e(S1)4t−e−(a26′)4t+ G260e−(a26′)4t* | 334 |
| $T_{24}^0 e^{(R_1)^{(4)}t} \le T_{24}(t) \le T_{24}^0 e^{((R_1)^{(4)}+(r_{24})^{(4)})t}$ | |
| $\frac{1}{(\mu_1)^{(4)}} T_{24}^0 e^{(R_1)^{(4)}t} \le T_{24}(t) \le \frac{1}{(\mu_2)^{(4)}} T_{24}^0 e^{((R_1)^{(4)}+(r_{24})^{(4)})t}$ | 335 |
| $\frac{(b_{26})^{(4)} T_{24}^0}{(\mu_1)^{(4)}((R_1)^{(4)}-(b_{26}')^{(4)})}\left[e^{(R_1)^{(4)}t} - e^{-(b_{26}')^{(4)}t}\right] + T_{26}^0 e^{-(b_{26}')^{(4)}t} \le T_{26}(t) \le$<br>$\frac{(a_{26})^{(4)} T_{24}^0}{(\mu_2)^{(4)}((R_1)^{(4)}+(r_{24})^{(4)}+(R_2)^{(4)})}\left[e^{((R_1)^{(4)}+(r_{24})^{(4)})t} - e^{-(R_2)^{(4)}t}\right] + T_{26}^0 e^{-(R_2)^{(4)}t}$ | 336 |
| **Definition of** $(S_1)^{(4)}, (S_2)^{(4)}, (R_1)^{(4)}, (R_2)^{(4)}$:-<br>Where $(S_1)^{(4)} = (a_{24})^{(4)}(m_2)^{(4)} - (a_{24}')^{(4)}$<br>$(S_2)^{(4)} = (a_{26})^{(4)} - (p_{26})^{(4)}$<br>$(R_1)^{(4)} = (b_{24})^{(4)}(\mu_2)^{(4)} - (b_{24}')^{(4)}$<br>$(R_2)^{(4)} = (b_{26}')^{(4)} - (r_{26})^{(4)}$ | 337 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ :<br>$(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ four constants satisfying<br>$-(\sigma_2)^{(5)} \le -(a_{28}')^{(5)} + (a_{29}')^{(5)} - (a_{28}'')^{(5)}(T_{29}, t) + (a_{29}'')^{(5)}(T_{29}, t) \le -(\sigma_1)^{(5)}$<br>$-(\tau_2)^{(5)} \le -(b_{28}')^{(5)} + (b_{29}')^{(5)} - (b_{28}'')^{(5)}((G_{31}), t) - (b_{29}'')^{(5)}((G_{31}), t) \le -(\tau_1)^{(5)}$ | 338 |
| **Definition of** $(\nu_1)^{(5)}, (\nu_2)^{(5)}, (u_1)^{(5)}, (u_2)^{(5)}, \nu^{(5)}, u^{(5)}$ :<br>By $(\nu_1)^{(5)} > 0$, $(\nu_2)^{(5)} < 0$ and respectively $(u_1)^{(5)} > 0$, $(u_2)^{(5)} < 0$ the roots of the equations $(a_{29})^{(5)}(\nu^{(5)})^2 + (\sigma_1)^{(5)}\nu^{(5)} - (a_{28})^{(5)} = 0$ | 339 |

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| and $(b_{29})^{(5)}\left(u^{(5)}\right)^2 + (\tau_1)^{(5)}u^{(5)} - (b_{28})^{(5)} = 0$ and | |
| **Definition of** $(\bar{v}_1)^{(5)},,(\bar{v}_2)^{(5)},(\bar{u}_1)^{(5)},(\bar{u}_2)^{(5)}$ :<br><br>By $(\bar{v}_1)^{(5)} > 0$ , $(\bar{v}_2)^{(5)} < 0$ and respectively $(\bar{u}_1)^{(5)} > 0$ , $(\bar{u}_2)^{(5)} < 0$ the roots of the equations $(a_{29})^{(5)}\left(v^{(5)}\right)^2 + (\sigma_2)^{(5)}v^{(5)} - (a_{28})^{(5)} = 0$ and $(b_{29})^{(5)}\left(u^{(5)}\right)^2 + (\tau_2)^{(5)}u^{(5)} - (b_{28})^{(5)} = 0$<br>**Definition of** $(m_1)^{(5)}$ , $(m_2)^{(5)}$ , $(\mu_1)^{(5)},(\mu_2)^{(5)},(v_0)^{(5)}$ :-<br><br>If we define $(m_1)^{(5)}$ , $(m_2)^{(5)}$ , $(\mu_1)^{(5)},(\mu_2)^{(5)}$ by<br><br>$(m_2)^{(5)} = (v_0)^{(5)}, (m_1)^{(5)} = (v_1)^{(5)}$, ***if*** $(v_0)^{(5)} < (v_1)^{(5)}$<br><br>$(m_2)^{(5)} = (v_1)^{(5)}, (m_1)^{(5)} = (\bar{v}_1)^{(5)}$ , ***if*** $(v_1)^{(5)} < (v_0)^{(5)} < (\bar{v}_1)^{(5)}$,<br>and $\boxed{(v_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}}$<br><br>$(m_2)^{(5)} = (v_1)^{(5)}, (m_1)^{(5)} = (v_0)^{(5)}$, ***if*** $(\bar{v}_1)^{(5)} < (v_0)^{(5)}$ | 340 |
| and analogously<br><br>$(\mu_2)^{(5)} = (u_0)^{(5)}, (\mu_1)^{(5)} = (u_1)^{(5)}$, ***if*** $(u_0)^{(5)} < (u_1)^{(5)}$<br><br>$(\mu_2)^{(5)} = (u_1)^{(5)}, (\mu_1)^{(5)} = (\bar{u}_1)^{(5)}$ , ***if*** $(u_1)^{(5)} < (u_0)^{(5)} < (\bar{u}_1)^{(5)}$,<br>and $\boxed{(u_0)^{(5)} = \frac{T_{28}^0}{T_{29}^0}}$<br><br>$(\mu_2)^{(5)} = (u_1)^{(5)}, (\mu_1)^{(5)} = (u_0)^{(5)}$, ***if*** $(\bar{u}_1)^{(5)} < (u_0)^{(5)}$ where $(u_1)^{(5)}, (\bar{u}_1)^{(5)}$ | 341 |
| Then the solution of global equations satisfies the inequalities<br><br>$G_{28}^0 e^{((S_1)^{(5)} - (p_{28})^{(5)})t} \le G_{28}(t) \le G_{28}^0 e^{(S_1)^{(5)}t}$<br><br>where $(p_i)^{(5)}$ is defined by equation | 342 |
| $\frac{1}{(m_5)^{(5)}} G_{28}^0 e^{((S_1)^{(5)} - (p_{28})^{(5)})t} \le G_{29}(t) \le \frac{1}{(m_2)^{(5)}} G_{28}^0 e^{(S_1)^{(5)}t}$ | 343 |
| $\left(\frac{(a_{30})^{(5)} G_{28}^0}{(m_1)^{(5)}((S_1)^{(5)} - (p_{28})^{(5)} - (S_2)^{(5)})}\left[e^{((S_1)^{(5)} - (p_{28})^{(5)})t} - e^{-(S_2)^{(5)}t}\right] + G_{30}^0 e^{-(S_2)^{(5)}t} \le G_{30}(t) \le$<br>*(a30)5G280(m2)5(S1)5−(a30′)5e(S1)5t−e−(a30′)5t+ G300e−(a30′)5t* | 344 |
| $\boxed{T_{28}^0 e^{(R_1)^{(5)}t} \le T_{28}(t) \le T_{28}^0 e^{((R_1)^{(5)} + (r_{28})^{(5)})t}}$ | 345 |
| $\frac{1}{(\mu_1)^{(5)}} T_{28}^0 e^{(R_1)^{(5)}t} \le T_{28}(t) \le \frac{1}{(\mu_2)^{(5)}} T_{28}^0 e^{((R_1)^{(5)} + (r_{28})^{(5)})t}$ | 346 |

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| $\frac{(b_{30})^{(5)}T_{28}^{0}}{(\mu_1)^{(5)}((R_1)^{(5)}-(b_{30}^{'})^{(5)})}\left[e^{(R_1)^{(5)}t}-e^{-(b_{30}^{'})^{(5)}t}\right]+T_{30}^{0}e^{-(b_{30}^{'})^{(5)}t}\le T_{30}(t)\le$ $\frac{(a_{30})^{(5)}T_{28}^{0}}{(\mu_2)^{(5)}((R_1)^{(5)}+(r_{28})^{(5)}+(R_2)^{(5)})}\left[e^{((R_1)^{(5)}+(r_{28})^{(5)})t}-e^{-(R_2)^{(5)}t}\right]+T_{30}^{0}e^{-(R_2)^{(5)}t}$ | 347 |
| **Definition of** $(S_1)^{(5)},(S_2)^{(5)},(R_1)^{(5)},(R_2)^{(5)}$:- Where $(S_1)^{(5)}=(a_{28})^{(5)}(m_2)^{(5)}-(a_{28}^{'})^{(5)}$ $(S_2)^{(5)}=(a_{30})^{(5)}-(p_{30})^{(5)}$ $(R_1)^{(5)}=(b_{28})^{(5)}(\mu_2)^{(5)}-(b_{28}^{'})^{(5)}$ $(R_2)^{(5)}=(b_{30}^{'})^{(5)}-(r_{30})^{(5)}$ | 348 |
| **Behavior of the solutions of equation** **Theorem 2:** If we denote and define | 349 |
| **Definition of** $(\sigma_1)^{(6)},(\sigma_2)^{(6)},(\tau_1)^{(6)},(\tau_2)^{(6)}$ : $(\sigma_1)^{(6)},(\sigma_2)^{(6)},(\tau_1)^{(6)},(\tau_2)^{(6)}$ four constants satisfying $-(\sigma_2)^{(6)}\le-(a_{32}^{'})^{(6)}+(a_{33}^{'})^{(6)}-(a_{32}^{''})^{(6)}(T_{33},t)+(a_{33}^{''})^{(6)}(T_{33},t)\le-(\sigma_1)^{(6)}$ $-(\tau_2)^{(6)}\le-(b_{32}^{'})^{(6)}+(b_{33}^{'})^{(6)}-(b_{32}^{''})^{(6)}\big((G_{35}),t\big)-(b_{33}^{''})^{(6)}\big((G_{35}),t\big)\le-(\tau_1)^{(6)}$ | 350 |
| **Definition of** $(\bar{\nu}_1)^{(6)},,(\bar{\nu}_2)^{(6)},(\bar{u}_1)^{(6)},(\bar{u}_2)^{(6)}$ : By $(\bar{\nu}_1)^{(6)}>0$ , $(\bar{\nu}_2)^{(6)}<0$ and respectively $(\bar{u}_1)^{(6)}>0$ , $(\bar{u}_2)^{(6)}<0$ the roots of the equations $(a_{33})^{(6)}\big(\nu^{(6)}\big)^2+(\sigma_2)^{(6)}\nu^{(6)}-(a_{32})^{(6)}=0$ and $(b_{33})^{(6)}\big(u^{(6)}\big)^2+(\tau_2)^{(6)}u^{(6)}-(b_{32})^{(6)}=0$ and **Definition of** $(m_1)^{(6)},(m_2)^{(6)},(\mu_1)^{(6)},(\mu_2)^{(6)},(\nu_0)^{(6)}$ :- If we define $(m_1)^{(6)},(m_2)^{(6)},(\mu_1)^{(6)},(\mu_2)^{(6)}$ by $(m_2)^{(6)}=(\nu_0)^{(6)},(m_1)^{(6)}=(\nu_1)^{(6)},$ ***if*** $(\nu_0)^{(6)}<(\nu_1)^{(6)}$ $(m_2)^{(6)}=(\nu_1)^{(6)},(m_1)^{(6)}=(\bar{\nu}_6)^{(6)}$ , ***if*** $(\nu_1)^{(6)}<(\nu_0)^{(6)}<(\bar{\nu}_1)^{(6)},$ and $\boxed{(\nu_0)^{(6)}=\frac{G_{32}^{0}}{G_{33}^{0}}}$ $(m_2)^{(6)}=(\nu_1)^{(6)},(m_1)^{(6)}=(\nu_0)^{(6)},$ ***if*** $(\bar{\nu}_1)^{(6)}<(\nu_0)^{(6)}$ | 351 |

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| and analogously<br>$(\mu_2)^{(6)} = (u_0)^{(6)}, (\mu_1)^{(6)} = (u_1)^{(6)},\ \boldsymbol{if}\ (u_0)^{(6)} < (u_1)^{(6)}$<br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (\bar{u}_1)^{(6)}, \boldsymbol{if}\ (u_1)^{(6)} < (u_0)^{(6)} < (\bar{u}_1)^{(6)},$<br>and $\boxed{(u_0)^{(6)} = \frac{T_{32}^0}{T_{33}^0}}$<br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (u_0)^{(6)}, \boldsymbol{if}\ (\bar{u}_1)^{(6)} < (u_0)^{(6)}$ where $(u_1)^{(6)}, (\bar{u}_1)^{(6)}$ | 352 |
| Then the solution of global equations satisfies the inequalities<br>$G_{32}^0 e^{((S_1)^{(6)} - (p_{32})^{(6)})t} \le G_{32}(t) \le G_{32}^0 e^{(S_1)^{(6)} t}$<br>where $(p_i)^{(6)}$ is defined by equation | 353 |
| $\frac{1}{(m_1)^{(6)}} G_{32}^0 e^{((S_1)^{(6)} - (p_{32})^{(6)})t} \le G_{33}(t) \le \frac{1}{(m_2)^{(6)}} G_{32}^0 e^{(S_1)^{(6)} t}$ | 354 |
| $\left(\frac{(a_{34})^{(6)} G_{32}^0}{(m_1)^{(6)}((S_1)^{(6)} - (p_{32})^{(6)} - (S_2)^{(6)})}\left[e^{((S_1)^{(6)} - (p_{32})^{(6)})t} - e^{-(S_2)^{(6)} t}\right] + G_{34}^0 e^{-(S_2)^{(6)} t} \le G_{34}(t) \le\right.$<br>*(a34)6G320(m2)6(S1)6−(a34′)6e(S1)6t−e−(a34′)6t+ G340e−(a34′)6t* | 355 |
| $\boxed{T_{32}^0 e^{(R_1)^{(6)} t} \le T_{32}(t) \le T_{32}^0 e^{((R_1)^{(6)} + (r_{32})^{(6)})t}}$ | 356 |
| $\frac{1}{(\mu_1)^{(6)}} T_{32}^0 e^{(R_1)^{(6)} t} \le T_{32}(t) \le \frac{1}{(\mu_2)^{(6)}} T_{32}^0 e^{((R_1)^{(6)} + (r_{32})^{(6)})t}$ | 357 |
| $\frac{(b_{34})^{(6)} T_{32}^0}{(\mu_1)^{(6)}((R_1)^{(6)} - (b_{34}')^{(6)})}\left[e^{(R_1)^{(6)} t} - e^{-(b_{34}')^{(6)} t}\right] + T_{34}^0 e^{-(b_{34}')^{(6)} t} \le T_{34}(t) \le$<br>$\frac{(a_{34})^{(6)} T_{32}^0}{(\mu_2)^{(6)}((R_1)^{(6)} + (r_{32})^{(6)} + (R_2)^{(6)})}\left[e^{((R_1)^{(6)} + (r_{32})^{(6)})t} - e^{-(R_2)^{(6)} t}\right] + T_{34}^0 e^{-(R_2)^{(6)} t}$ | 358 |
| **Definition of** $(S_1)^{(6)}, (S_2)^{(6)}, (R_1)^{(6)}, (R_2)^{(6)}$:-<br>Where $(S_1)^{(6)} = (a_{32})^{(6)}(m_2)^{(6)} - (a_{32}')^{(6)}$<br>$(S_2)^{(6)} = (a_{34})^{(6)} - (p_{34})^{(6)}$<br>$(R_1)^{(6)} = (b_{32})^{(6)}(\mu_2)^{(6)} - (b_{32}')^{(6)}$<br>$(R_2)^{(6)} = (b_{34}')^{(6)} - (r_{34})^{(6)}$ | 359 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(7)}, (\sigma_2)^{(7)}, (\tau_1)^{(7)}, (\tau_2)^{(7)}$ : | |

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| $(\sigma_1)^{(7)}, (\sigma_2)^{(7)}, (\tau_1)^{(7)}, (\tau_2)^{(7)}$ four constants satisfying<br><br>$-(\sigma_2)^{(7)} \le -(a'_{36})^{(7)} + (a'_{37})^{(7)} - (a''_{36})^{(7)}(T_{37}, t) + (a''_{37})^{(7)}(T_{37}, t) \le -(\sigma_1)^{(7)}$<br><br>$-(\tau_2)^{(7)} \le -(b'_{36})^{(7)} + (b'_{37})^{(7)} - (b''_{36})^{(7)}((G_{39}), t) - (b''_{37})^{(7)}((G_{39}), t) \le -(\tau_1)^{(7)}$ | |
| **Definition of** $(\nu_1)^{(7)}, (\nu_2)^{(7)}, (u_1)^{(7)}, (u_2)^{(7)}, \nu^{(7)}, u^{(7)}$ :<br><br>By $(\nu_1)^{(7)} > 0, (\nu_2)^{(7)} < 0$ and respectively $(u_1)^{(7)} > 0, (u_2)^{(7)} < 0$ the roots of the equations $(a_{37})^{(7)}(\nu^{(7)})^2 + (\sigma_1)^{(7)}\nu^{(7)} - (a_{36})^{(7)} = 0$<br>and $(b_{37})^{(7)}(u^{(7)})^2 + (\tau_1)^{(7)}u^{(7)} - (b_{36})^{(7)} = 0$ and | 361 |
| **Definition of** $(\bar{\nu}_1)^{(7)},, (\bar{\nu}_2)^{(7)}, (\bar{u}_1)^{(7)}, (\bar{u}_2)^{(7)}$ :<br><br>By $(\bar{\nu}_1)^{(7)} > 0, (\bar{\nu}_2)^{(7)} < 0$ and respectively $(\bar{u}_1)^{(7)} > 0, (\bar{u}_2)^{(7)} < 0$ the roots of the equations $(a_{37})^{(7)}(\nu^{(7)})^2 + (\sigma_2)^{(7)}\nu^{(7)} - (a_{36})^{(7)} = 0$<br>and $(b_{37})^{(7)}(u^{(7)})^2 + (\tau_2)^{(7)}u^{(7)} - (b_{36})^{(7)} = 0$<br>**Definition of** $(m_1)^{(7)}, (m_2)^{(7)}, (\mu_1)^{(7)}, (\mu_2)^{(7)}, (\nu_0)^{(7)}$ :-<br><br>If we define $(m_1)^{(7)}, (m_2)^{(7)}, (\mu_1)^{(7)}, (\mu_2)^{(7)}$ by<br><br>$(m_2)^{(7)} = (\nu_0)^{(7)}, (m_1)^{(7)} = (\nu_1)^{(7)}$, ***if*** $(\nu_0)^{(7)} < (\nu_1)^{(7)}$<br><br>$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\bar{\nu}_1)^{(7)}$, ***if*** $(\nu_1)^{(7)} < (\nu_0)^{(7)} < (\bar{\nu}_1)^{(7)}$,<br>and $\boxed{(\nu_0)^{(7)} = \frac{G^0_{36}}{G^0_{37}}}$<br><br>$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\nu_0)^{(7)}$, ***if*** $(\bar{\nu}_1)^{(7)} < (\nu_0)^{(7)}$ | 362 |
| and analogously<br><br>$(\mu_2)^{(7)} = (u_0)^{(7)}, (\mu_1)^{(7)} = (u_1)^{(7)}$, ***if*** $(u_0)^{(7)} < (u_1)^{(7)}$<br><br>$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (\bar{u}_1)^{(7)}$, ***if*** $(u_1)^{(7)} < (u_0)^{(7)} < (\bar{u}_1)^{(7)}$,<br>and $\boxed{(u_0)^{(7)} = \frac{T^0_{36}}{T^0_{37}}}$<br><br>$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (u_0)^{(7)}$, ***if*** $(\bar{u}_1)^{(7)} < (u_0)^{(7)}$ where $(u_1)^{(7)}, (\bar{u}_1)^{(7)}$ | 363 |
| Then the solution of global equations satisfies the inequalities | 364 |

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|---|---|
| $G_{36}^{0}e^{((S_1)^{(7)}-(p_{36})^{(7)})t} \le G_{36}(t) \le G_{36}^{0}e^{(S_1)^{(7)}t}$<br><br>where $(p_i)^{(7)}$ is defined by equation | |
| $\frac{1}{(m_7)^{(7)}} G_{36}^{0}e^{((S_1)^{(7)}-(p_{36})^{(7)})t} \le G_{37}(t) \le \frac{1}{(m_2)^{(7)}} G_{36}^{0}e^{(S_1)^{(7)}t}$ | 365 |
| $(\frac{(a_{38})^{(7)}G_{36}^{0}}{(m_1)^{(7)}((S_1)^{(7)}-(p_{36})^{(7)}-(S_2)^{(7)})}\left[e^{((S_1)^{(7)}-(p_{36})^{(7)})t} - e^{-(S_2)^{(7)}t}\right] + G_{38}^{0}e^{-(S_2)^{(7)}t} \le G_{38}(t) \le$<br>$\frac{(a_{38})^{(7)}G_{36}^{0}}{(m_2)^{(7)}((S_1)^{(7)}-(a_{38}^{'})^{(7)})}[e^{(S_1)^{(7)}t} - e^{-(a_{38}^{'})^{(7)}t}] + G_{38}^{0}e^{-(a_{38}^{'})^{(7)}t})$ | 366 |
| $\boxed{T_{36}^{0}e^{(R_1)^{(7)}t} \le T_{36}(t) \le T_{36}^{0}e^{((R_1)^{(7)}+(r_{36})^{(7)})t}}$ | 367 |
| $\frac{1}{(\mu_1)^{(7)}} T_{36}^{0}e^{(R_1)^{(7)}t} \le T_{36}(t) \le \frac{1}{(\mu_2)^{(7)}} T_{36}^{0}e^{((R_1)^{(7)}+(r_{36})^{(7)})t}$ | 368 |
| $\frac{(b_{38})^{(7)}T_{36}^{0}}{(\mu_1)^{(7)}((R_1)^{(7)}-(b_{38}^{'})^{(7)})}\left[e^{(R_1)^{(7)}t} - e^{-(b_{38}^{'})^{(7)}t}\right] + T_{38}^{0}e^{-(b_{38}^{'})^{(7)}t} \le T_{38}(t) \le$<br><br>$\frac{(a_{38})^{(7)}T_{36}^{0}}{(\mu_2)^{(7)}((R_1)^{(7)}+(r_{36})^{(7)}+(R_2)^{(7)})}\left[e^{((R_1)^{(7)}+(r_{36})^{(7)})t} - e^{-(R_2)^{(7)}t}\right] + T_{38}^{0}e^{-(R_2)^{(7)}t}$ | 369 |
| **Definition of** $(S_1)^{(7)}, (S_2)^{(7)}, (R_1)^{(7)}, (R_2)^{(7)}$:-<br><br>Where $(S_1)^{(7)} = (a_{36})^{(7)}(m_2)^{(7)} - (a_{36}^{'})^{(7)}$<br><br>$(S_2)^{(7)} = (a_{38})^{(7)} - (p_{38})^{(7)}$<br><br>$(R_1)^{(7)} = (b_{36})^{(7)}(\mu_2)^{(7)} - (b_{36}^{'})^{(7)}$<br><br>$(R_2)^{(7)} = (b_{38}^{'})^{(7)} - (r_{38})^{(7)}$ | 370 |
| **Behavior of the solutions of equation**<br><br>**Theorem 2:** If we denote and define<br><br>**Definition of** $(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ :<br><br>$(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ four constants satisfying<br><br>$-(\sigma_2)^{(8)} \le -(a_{40}^{'})^{(8)} + (a_{41}^{'})^{(8)} - (a_{40}^{''})^{(8)}(T_{41}, t) + (a_{41}^{''})^{(8)}(T_{41}, t) \le -(\sigma_1)^{(8)}$<br><br>$-(\tau_2)^{(8)} \le -(b_{40}^{'})^{(8)} + (b_{41}^{'})^{(8)} - (b_{40}^{''})^{(8)}((G_{43}), t) - (b_{41}^{''})^{(8)}((G_{43}), t) \le -(\tau_1)^{(8)}$ | 371 |
| **Definition of** $(v_1)^{(8)}, (v_2)^{(8)}, (u_1)^{(8)}, (u_2)^{(8)}, v^{(8)}, u^{(8)}$ : | 372 |

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| By $(\nu_1)^{(8)} > 0\,, (\nu_2)^{(8)} < 0$ and respectively $(u_1)^{(8)} > 0\,, (u_2)^{(8)} < 0$ the roots of the equations $(a_{41})^{(8)}\left(\nu^{(8)}\right)^2 + (\sigma_1)^{(8)}\nu^{(8)} - (a_{40})^{(8)} = 0$ and $(b_{41})^{(8)}\left(u^{(8)}\right)^2 + (\tau_1)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$ and | |
| **Definition of** $(\bar{\nu}_1)^{(8)},, (\bar{\nu}_2)^{(8)}, (\bar{u}_1)^{(8)}, (\bar{u}_2)^{(8)}$ :<br><br>By $(\bar{\nu}_1)^{(8)} > 0\,, (\bar{\nu}_2)^{(8)} < 0$ and respectively $(\bar{u}_1)^{(8)} > 0\,, (\bar{u}_2)^{(8)} < 0$ the<br><br>roots of the equations $(a_{41})^{(8)}\left(\nu^{(8)}\right)^2 + (\sigma_2)^{(8)}\nu^{(8)} - (a_{40})^{(8)} = 0$<br><br>and $(b_{41})^{(8)}\left(u^{(8)}\right)^2 + (\tau_2)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$<br><br>**Definition of** $(m_1)^{(8)}\,, (m_2)^{(8)}\,, (\mu_1)^{(8)}, (\mu_2)^{(8)}, (\nu_0)^{(8)}$ :-<br><br>If we define $(m_1)^{(8)}\,, (m_2)^{(8)}\,, (\mu_1)^{(8)}, (\mu_2)^{(8)}$ by<br><br>$(m_2)^{(8)} = (\nu_0)^{(8)}, (m_1)^{(8)} = (\nu_1)^{(8)},$ ***if*** $(\nu_0)^{(8)} < (\nu_1)^{(8)}$<br><br>$(m_2)^{(8)} = (\nu_1)^{(8)}, (m_1)^{(8)} = (\bar{\nu}_1)^{(8)}$ , ***if*** $(\nu_1)^{(8)} < (\nu_0)^{(8)} < (\bar{\nu}_1)^{(8)},$<br>and $\boxed{(\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}}$<br><br>$(m_2)^{(8)} = (\nu_1)^{(8)}, (m_1)^{(8)} = (\nu_0)^{(8)},$ ***if*** $(\bar{\nu}_1)^{(8)} < (\nu_0)^{(8)}$ | |
| and analogously<br><br>$(\mu_2)^{(8)} = (u_0)^{(8)}, (\mu_1)^{(8)} = (u_1)^{(8)},$ ***if*** $(u_0)^{(8)} < (u_1)^{(8)}$<br><br>$(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (\bar{u}_1)^{(8)}$ , ***if*** $(u_1)^{(8)} < (u_0)^{(8)} < (\bar{u}_1)^{(8)},$<br>and $\boxed{(u_0)^{(8)} = \frac{T_{40}^0}{T_{41}^0}}$<br><br>$(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (u_0)^{(8)},$ ***if*** $(\bar{u}_1)^{(8)} < (u_0)^{(8)}$ where $(u_1)^{(8)}, (\bar{u}_1)^{(8)}$ | 374 |
| Then the solution of global equations satisfies the inequalities<br><br>$G_{40}^0 e^{((S_1)^{(8)} - (p_{40})^{(8)})t} \le G_{40}(t) \le G_{40}^0 e^{(S_1)^{(8)}t}$<br><br>where $(p_i)^{(8)}$ is defined by equation | 375 |
| $\frac{1}{(m_1)^{(8)}} G_{40}^0 e^{((S_1)^{(8)} - (p_{40})^{(8)})t} \le G_{41}(t) \le \frac{1}{(m_2)^{(8)}} G_{40}^0 e^{(S_1)^{(8)}t}$ | 376 |

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| $\left(\frac{(a_{42})^{(8)}G_{40}^{0}}{(m_1)^{(8)}((S_1)^{(8)}-(p_{40})^{(8)}-(S_2)^{(8)})}\left[e^{((S_1)^{(8)}-(p_{40})^{(8)})t}-e^{-(S_2)^{(8)}t}\right]+G_{42}^{0}e^{-(S_2)^{(8)}t}\leq G_{42}(t)\leq \frac{(a_{42})^{(8)}G_{40}^{0}}{(m_2)^{(8)}((S_1)^{(8)}-(a_{42}^{'})^{(8)})}[e^{(S_1)^{(8)}t}-e^{-(a_{42}^{'})^{(8)}t}]+G_{42}^{0}e^{-(a_{42}^{'})^{(8)}t}\right)$ | 377 |
| $\boxed{T_{40}^{0}e^{(R_1)^{(8)}t}\leq T_{40}(t)\leq T_{40}^{0}e^{((R_1)^{(8)}+(r_{40})^{(8)})t}}$ | 378 |
| $\frac{1}{(\mu_1)^{(8)}}T_{40}^{0}e^{(R_1)^{(8)}t}\leq T_{40}(t)\leq \frac{1}{(\mu_2)^{(8)}}T_{40}^{0}e^{((R_1)^{(8)}+(r_{40})^{(8)})t}$ | 379 |
| $\frac{(b_{42})^{(8)}T_{40}^{0}}{(\mu_1)^{(8)}((R_1)^{(8)}-(b_{42}^{'})^{(8)})}\left[e^{(R_1)^{(8)}t}-e^{-(b_{42}^{'})^{(8)}t}\right]+T_{42}^{0}e^{-(b_{42}^{'})^{(8)}t}\leq T_{42}(t)\leq$ $\frac{(a_{42})^{(8)}T_{40}^{0}}{(\mu_2)^{(8)}((R_1)^{(8)}+(r_{40})^{(8)}+(R_2)^{(8)})}\left[e^{((R_1)^{(8)}+(r_{40})^{(8)})t}-e^{-(R_2)^{(8)}t}\right]+T_{42}^{0}e^{-(R_2)^{(8)}t}$ | 380 |
| **Definition of** $(S_1)^{(8)},(S_2)^{(8)},(R_1)^{(8)},(R_2)^{(8)}$:-<br>Where $(S_1)^{(8)}=(a_{40})^{(8)}(m_2)^{(8)}-(a_{40}^{'})^{(8)}$<br>$(S_2)^{(8)}=(a_{42})^{(8)}-(p_{42})^{(8)}$<br>$(R_1)^{(8)}=(b_{40})^{(8)}(\mu_2)^{(8)}-(b_{40}^{'})^{(8)}$<br>$(R_2)^{(8)}=(b_{42}^{'})^{(8)}-(r_{42})^{(8)}$ | 381 |
| **Behavior of the solutions of equation 37 to 92**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(9)},(\sigma_2)^{(9)},(\tau_1)^{(9)},(\tau_2)^{(9)}$ :<br>$(\sigma_1)^{(9)},(\sigma_2)^{(9)},(\tau_1)^{(9)},(\tau_2)^{(9)}$ four constants satisfying<br>$-(\sigma_2)^{(9)}\leq -(a_{44}^{'})^{(9)}+(a_{45}^{'})^{(9)}-(a_{44}^{''})^{(9)}(T_{45},t)+(a_{45}^{''})^{(9)}(T_{45},t)\leq -(\sigma_1)^{(9)}$<br>$-(\tau_2)^{(9)}\leq -(b_{44}^{'})^{(9)}+(b_{45}^{'})^{(9)}-(b_{44}^{''})^{(9)}((G_{47}),t)-(b_{45}^{''})^{(9)}((G_{47}),t)\leq -(\tau_1)^{(9)}$ | 382 |
| **Definition of** $(\nu_1)^{(9)},(\nu_2)^{(9)},(u_1)^{(9)},(u_2)^{(9)},\nu^{(9)},u^{(9)}$ :<br>By $(\nu_1)^{(9)}>0$ , $(\nu_2)^{(9)}<0$ and respectively $(u_1)^{(9)}>0$ , $(u_2)^{(9)}<0$ the roots of the equations $(a_{45})^{(9)}(\nu^{(9)})^2+(\sigma_1)^{(9)}\nu^{(9)}-(a_{44})^{(9)}=0$<br>and $(b_{45})^{(9)}(u^{(9)})^2+(\tau_1)^{(9)}u^{(9)}-(b_{44})^{(9)}=0$ and | |
| **Definition of** $(\bar{\nu}_1)^{(9)},,(\bar{\nu}_2)^{(9)},(\bar{u}_1)^{(9)},(\bar{u}_2)^{(9)}$ :<br>By $(\bar{\nu}_1)^{(9)}>0$ , $(\bar{\nu}_2)^{(9)}<0$ and respectively $(\bar{u}_1)^{(9)}>0$ , $(\bar{u}_2)^{(9)}<0$ the roots of the equations $(a_{45})^{(9)}(\nu^{(9)})^2+(\sigma_2)^{(9)}\nu^{(9)}-(a_{44})^{(9)}=0$ | |

and $(b_{45})^{(9)}\left(u^{(9)}\right)^2+(\tau_2)^{(9)}u^{(9)}-(b_{44})^{(9)}=0$

**Definition of** $(m_1)^{(9)},(m_2)^{(9)},(\mu_1)^{(9)},(\mu_2)^{(9)},(\nu_0)^{(9)}$ :-

If we define $(m_1)^{(9)},(m_2)^{(9)},(\mu_1)^{(9)},(\mu_2)^{(9)}$ by

$(m_2)^{(9)}=(\nu_0)^{(9)},(m_1)^{(9)}=(\nu_1)^{(9)},$ ***if*** $(\nu_0)^{(9)}<(\nu_1)^{(9)}$

$(m_2)^{(9)}=(\nu_1)^{(9)},(m_1)^{(9)}=(\bar{\nu}_1)^{(9)}$ , ***if*** $(\nu_1)^{(9)}<(\nu_0)^{(9)}<(\bar{\nu}_1)^{(9)}$,

and $\boxed{(\nu_0)^{(9)}=\frac{G_{44}^0}{G_{45}^0}}$

$(m_2)^{(9)}=(\nu_1)^{(9)},(m_1)^{(9)}=(\nu_0)^{(9)},$ ***if*** $(\bar{\nu}_1)^{(9)}<(\nu_0)^{(9)}$

and analogously

$(\mu_2)^{(9)}=(u_0)^{(9)},(\mu_1)^{(9)}=(u_1)^{(9)},$ ***if*** $(u_0)^{(9)}<(u_1)^{(9)}$

$(\mu_2)^{(9)}=(u_1)^{(9)},(\mu_1)^{(9)}=(\bar{u}_1)^{(9)}$ , ***if*** $(u_1)^{(9)}<(u_0)^{(9)}<(\bar{u}_1)^{(9)}$,

and $\boxed{(u_0)^{(9)}=\frac{T_{44}^0}{T_{45}^0}}$

$(\mu_2)^{(9)}=(u_1)^{(9)},(\mu_1)^{(9)}=(u_0)^{(9)},$ ***if*** $(\bar{u}_1)^{(9)}<(u_0)^{(9)}$ where $(u_1)^{(9)},(\bar{u}_1)^{(9)}$ are defined by 59 and 69 respectively

Then the solution of 19,20,21,22,23 and 24 satisfies the inequalities

$$G_{44}^0e^{((S_1)^{(9)}-(p_{44})^{(9)})t}\le G_{44}(t)\le G_{44}^0e^{(S_1)^{(9)}t}$$

where $(p_i)^{(9)}$ is defined by equation 45

$$\frac{1}{(m_9)^{(9)}}G_{44}^0e^{((S_1)^{(9)}-(p_{44})^{(9)})t}\le G_{45}(t)\le\frac{1}{(m_2)^{(9)}}G_{44}^0e^{(S_1)^{(9)}t}$$

$$\left(\frac{(a_{46})^{(9)}G_{44}^0}{(m_1)^{(9)}((S_1)^{(9)}-(p_{44})^{(9)}-(S_2)^{(9)})}\left[e^{((S_1)^{(9)}-(p_{44})^{(9)})t}-e^{-(S_2)^{(9)}t}\right]+G_{46}^0e^{-(S_2)^{(9)}t}\le G_{46}(t)\le\right.$$

$$\left.\frac{(a_{46})^{(9)}G_{44}^0}{(m_2)^{(9)}((S_1)^{(9)}-(a_{46}')^{(9)})}[e^{(S_1)^{(9)}t}-e^{-(a_{46}')^{(9)}t}]+\ G_{46}^0e^{-(a_{46}')^{(9)}t}\right)$$

$$\boxed{T_{44}^0e^{(R_1)^{(9)}t}\le T_{44}(t)\le T_{44}^0e^{((R_1)^{(9)}+(r_{44})^{(9)})t}}$$

$$\frac{1}{(\mu_1)^{(9)}}T_{44}^0e^{(R_1)^{(9)}t}\le T_{44}(t)\le\frac{1}{(\mu_2)^{(9)}}T_{44}^0e^{((R_1)^{(9)}+(r_{44})^{(9)})t}$$

$$\frac{(b_{46})^{(9)}T_{44}^0}{(\mu_1)^{(9)}((R_1)^{(9)}-(b_{46}')^{(9)})}\left[e^{(R_1)^{(9)}t}-e^{-(b_{46}')^{(9)}t}\right]+T_{46}^0e^{-(b_{46}')^{(9)}t}\le T_{46}(t)\le$$

$$\frac{(a_{46})^{(9)}T_{44}^0}{(\mu_2)^{(9)}((R_1)^{(9)}+(r_{44})^{(9)}+(R_2)^{(9)})}\left[e^{((R_1)^{(9)}+(r_{44})^{(9)})t}-e^{-(R_2)^{(9)}t}\right]+T_{46}^0e^{-(R_2)^{(9)}t}$$

**Definition of** $(S_1)^{(9)}, (S_2)^{(9)}, (R_1)^{(9)}, (R_2)^{(9)}$:-

Where $(S_1)^{(9)} = (a_{44})^{(9)}(m_2)^{(9)} - (a'_{44})^{(9)}$

$(S_2)^{(9)} = (a_{46})^{(9)} - (p_{46})^{(9)}$

$(R_1)^{(9)} = (b_{44})^{(9)}(\mu_2)^{(9)} - (b'_{44})^{(9)}$

$(R_2)^{(9)} = (b'_{46})^{(9)} - (r_{46})^{(9)}$

**Proof :** From global equations we obtain (383)

$$\frac{dv^{(1)}}{dt} = (a_{13})^{(1)} - \left((a'_{13})^{(1)} - (a'_{14})^{(1)} + (a''_{13})^{(1)}(T_{14}, t)\right) - (a''_{14})^{(1)}(T_{14}, t)v^{(1)} - (a_{14})^{(1)}v^{(1)}$$

**Definition of** $v^{(1)}$ :- $\boxed{v^{(1)} = \frac{G_{13}}{G_{14}}}$

It follows

$$-\left((a_{14})^{(1)}\left(v^{(1)}\right)^2 + (\sigma_2)^{(1)}v^{(1)} - (a_{13})^{(1)}\right) \le \frac{dv^{(1)}}{dt} \le -\left((a_{14})^{(1)}\left(v^{(1)}\right)^2 + (\sigma_1)^{(1)}v^{(1)} - (a_{13})^{(1)}\right)$$

From which one obtains

**Definition of** $(\bar{v}_1)^{(1)}, (v_0)^{(1)}$ :-

For $0 < \boxed{(v_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}} < (v_1)^{(1)} < (\bar{v}_1)^{(1)}$

$$v^{(1)}(t) \ge \frac{(v_1)^{(1)} + (C)^{(1)}(v_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((v_1)^{(1)} - (v_0)^{(1)}\right)t\right]}}{1 + (C)^{(1)}e^{\left[-(a_{14})^{(1)}\left((v_1)^{(1)} - (v_0)^{(1)}\right)t\right]}} \quad , \quad \boxed{(C)^{(1)} = \frac{(v_1)^{(1)} - (v_0)^{(1)}}{(v_0)^{(1)} - (v_2)^{(1)}}}$$

it follows $(v_0)^{(1)} \le v^{(1)}(t) \le (v_1)^{(1)}$

In the same manner , we get (384)

$$v^{(1)}(t) \le \frac{(\bar{v}_1)^{(1)} + (\bar{C})^{(1)}(\bar{v}_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{v}_1)^{(1)} - (\bar{v}_2)^{(1)}\right)t\right]}}{1 + (\bar{C})^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{v}_1)^{(1)} - (\bar{v}_2)^{(1)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(1)} = \frac{(\bar{v}_1)^{(1)} - (v_0)^{(1)}}{(v_0)^{(1)} - (\bar{v}_2)^{(1)}}}$$

From which we deduce $(v_0)^{(1)} \le v^{(1)}(t) \le (\bar{v}_1)^{(1)}$

If $0 < (v_1)^{(1)} < (v_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0} < (\bar{v}_1)^{(1)}$ we find like in the previous case, (385)

$$(v_1)^{(1)} \le \frac{(v_1)^{(1)} + (C)^{(1)}(v_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((v_1)^{(1)} - (v_2)^{(1)}\right)t\right]}}{1 + (C)^{(1)}e^{\left[-(a_{14})^{(1)}\left((v_1)^{(1)} - (v_2)^{(1)}\right)t\right]}} \le v^{(1)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(1)}+(\bar{C})^{(1)}(\bar{\nu}_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}}{1+(\bar{C})^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}} \leq (\bar{\nu}_1)^{(1)}$$

If $0 < (\nu_1)^{(1)} \leq (\bar{\nu}_1)^{(1)} \leq \boxed{(\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}}$ , we obtain 386

$$(\nu_1)^{(1)} \leq \nu^{(1)}(t) \leq \frac{(\bar{\nu}_1)^{(1)}+(\bar{C})^{(1)}(\bar{\nu}_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}}{1+(\bar{C})^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}} \leq (\nu_0)^{(1)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(1)}(t)$ :-

$$(m_2)^{(1)} \leq \nu^{(1)}(t) \leq (m_1)^{(1)}, \quad \boxed{\nu^{(1)}(t) = \frac{G_{13}(t)}{G_{14}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(1)}(t)$ :-

$$(\mu_2)^{(1)} \leq u^{(1)}(t) \leq (\mu_1)^{(1)}, \quad \boxed{u^{(1)}(t) = \frac{T_{13}(t)}{T_{14}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{13}^{''})^{(1)} = (a_{14}^{''})^{(1)}, then\ (\sigma_1)^{(1)} = (\sigma_2)^{(1)}$ and in this case $(\nu_1)^{(1)} = (\bar{\nu}_1)^{(1)}$ if in addition $(\nu_0)^{(1)} = (\nu_1)^{(1)}$ then $\nu^{(1)}(t) = (\nu_0)^{(1)}$ and as a consequence $G_{13}(t) = (\nu_0)^{(1)}G_{14}(t)$ this also defines $(\nu_0)^{(1)}$ for the special case

Analogously if $(b_{13}^{''})^{(1)} = (b_{14}^{''})^{(1)}, then\ (\tau_1)^{(1)} = (\tau_2)^{(1)}$ and then

$(u_1)^{(1)} = (\bar{u}_1)^{(1)}$if in addition $(u_0)^{(1)} = (u_1)^{(1)}$ then $T_{13}(t) = (u_0)^{(1)}T_{14}(t)$ This is an important consequence of the relation between $(\nu_1)^{(1)}$ and $(\bar{\nu}_1)^{(1)}$, and definition of $(u_0)^{(1)}$.

**Proof** : From global equations we obtain 387

$$\frac{d\nu^{(2)}}{dt} = (a_{16})^{(2)} - \left((a_{16}^{'})^{(2)} - (a_{17}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17}, t)\right) - (a_{17}^{''})^{(2)}(T_{17}, t)\nu^{(2)} - (a_{17})^{(2)}\nu^{(2)}$$

**Definition of** $\nu^{(2)}$ :- $\boxed{\nu^{(2)} = \frac{G_{16}}{G_{17}}}$ 388

It follows 389

$$-\left((a_{17})^{(2)}\left(\nu^{(2)}\right)^2 + (\sigma_2)^{(2)}\nu^{(2)} - (a_{16})^{(2)}\right) \leq \frac{d\nu^{(2)}}{dt} \leq -\left((a_{17})^{(2)}\left(\nu^{(2)}\right)^2 + (\sigma_1)^{(2)}\nu^{(2)} - (a_{16})^{(2)}\right)$$

From which one obtains 390

**Definition of** $(\bar{\nu}_1)^{(2)}, (\nu_0)^{(2)}$ :-

For $0 < (\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0} < (\nu_1)^{(2)} < (\bar{\nu}_1)^{(2)}$

$$\nu^{(2)}(t) \geq \frac{(\nu_1)^{(2)} + (C)^{(2)}(\nu_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_0)^{(2)}\right)t\right]}}{1 + (C)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_0)^{(2)}\right)t\right]}} \quad , \quad \boxed{(C)^{(2)} = \frac{(\nu_1)^{(2)} - (\nu_0)^{(2)}}{(\nu_0)^{(2)} - (\nu_2)^{(2)}}}$$

it follows $(\nu_0)^{(2)} \leq \nu^{(2)}(t) \leq (\nu_1)^{(2)}$

In the same manner , we get 391

$$\nu^{(2)}(t) \leq \frac{(\bar{\nu}_1)^{(2)} + (\bar{C})^{(2)}(\bar{\nu}_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}}{1 + (\bar{C})^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(2)} = \frac{(\bar{\nu}_1)^{(2)} - (\nu_0)^{(2)}}{(\nu_0)^{(2)} - (\bar{\nu}_2)^{(2)}}}$$

From which we deduce $(\nu_0)^{(2)} \leq \nu^{(2)}(t) \leq (\bar{\nu}_1)^{(2)}$ 392

If $0 < (\nu_1)^{(2)} < (\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0} < (\bar{\nu}_1)^{(2)}$ we find like in the previous case, 393

$$(\nu_1)^{(2)} \leq \frac{(\nu_1)^{(2)} + (C)^{(2)}(\nu_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_2)^{(2)}\right)t\right]}}{1 + (C)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\nu_1)^{(2)} - (\nu_2)^{(2)}\right)t\right]}} \leq \nu^{(2)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(2)} + (\bar{C})^{(2)}(\bar{\nu}_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}}{1 + (\bar{C})^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}} \leq (\bar{\nu}_1)^{(2)}$$

If $0 < (\nu_1)^{(2)} \leq (\bar{\nu}_1)^{(2)} \leq (\nu_0)^{(2)} = \frac{G_{16}^0}{G_{17}^0}$ , we obtain 394

$$(\nu_1)^{(2)} \leq \nu^{(2)}(t) \leq \frac{(\bar{\nu}_1)^{(2)} + (\bar{C})^{(2)}(\bar{\nu}_2)^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}}{1 + (\bar{C})^{(2)} e^{\left[-(a_{17})^{(2)}\left((\bar{\nu}_1)^{(2)} - (\bar{\nu}_2)^{(2)}\right)t\right]}} \leq (\nu_0)^{(2)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(2)}(t)$ :- 395

$$(m_2)^{(2)} \leq \nu^{(2)}(t) \leq (m_1)^{(2)}, \quad \boxed{\nu^{(2)}(t) = \frac{G_{16}(t)}{G_{17}(t)}}$$

In a completely analogous way, we obtain 396

**Definition of** $u^{(2)}(t)$ :-

$$(\mu_2)^{(2)} \leq u^{(2)}(t) \leq (\mu_1)^{(2)}, \quad \boxed{u^{(2)}(t) = \frac{T_{16}(t)}{T_{17}(t)}}$$

| | |
|---|---|
| Now, using this result and replacing it in global equations we get easily the result stated in the theorem. | |
| **Particular case :**<br><br>If $(a_{16}^{''})^{(2)} = (a_{17}^{''})^{(2)}, then\ (\sigma_1)^{(2)} = (\sigma_2)^{(2)}$ and in this case $(\nu_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ if in addition $(\nu_0)^{(2)} = (\nu_1)^{(2)}$ then $\nu^{(2)}(t) = (\nu_0)^{(2)}$ and as a consequence $G_{16}(t) = (\nu_0)^{(2)} G_{17}(t)$<br><br>Analogously if $(b_{16}^{''})^{(2)} = (b_{17}^{''})^{(2)}, then\ (\tau_1)^{(2)} = (\tau_2)^{(2)}$ and then<br><br>$(u_1)^{(2)} = (\bar{u}_1)^{(2)}$ if in addition $(u_0)^{(2)} = (u_1)^{(2)}$ then $T_{16}(t) = (u_0)^{(2)} T_{17}(t)$ This is an important consequence of the relation between $(\nu_1)^{(2)}$ and $(\bar{\nu}_1)^{(2)}$ | 397 |
| **Proof** : From global equations we obtain<br><br>$\frac{d\nu^{(3)}}{dt} = (a_{20})^{(3)} - \left((a_{20}^{'})^{(3)} - (a_{21}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21}, t)\right) - (a_{21}^{''})^{(3)}(T_{21}, t)\nu^{(3)} - (a_{21})^{(3)}\nu^{(3)}$ | 398 |
| **Definition of** $\nu^{(3)}$ :- $\boxed{\nu^{(3)} = \frac{G_{20}}{G_{21}}}$<br><br>It follows<br><br>$-\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_2)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right) \leq \frac{d\nu^{(3)}}{dt} \leq -\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_1)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right)$ | 399 |
| From which one obtains<br><br>For $0 < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\nu_1)^{(3)} < (\bar{\nu}_1)^{(3)}$<br><br>$\nu^{(3)}(t) \geq \frac{(\nu_1)^{(3)} + (C)^{(3)}(\nu_2)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)} - (\nu_0)^{(3)}\right)t\right]}}{1 + (C)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)} - (\nu_0)^{(3)}\right)t\right]}}$ , $\boxed{(C)^{(3)} = \frac{(\nu_1)^{(3)} - (\nu_0)^{(3)}}{(\nu_0)^{(3)} - (\nu_2)^{(3)}}}$<br><br>it follows $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\nu_1)^{(3)}$ | 400 |
| In the same manner , we get<br><br>$\nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)} + (\bar{C})^{(3)}(\bar{\nu}_2)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)} - (\bar{\nu}_2)^{(3)}\right)t\right]}}{1 + (C)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)} - (\bar{\nu}_2)^{(3)}\right)t\right]}}$ , $\boxed{(\bar{C})^{(3)} = \frac{(\bar{\nu}_1)^{(3)} - (\nu_0)^{(3)}}{(\nu_0)^{(3)} - (\bar{\nu}_2)^{(3)}}}$<br><br>**Definition of** $(\bar{\nu}_1)^{(3)}$ :-<br><br>From which we deduce $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\bar{\nu}_1)^{(3)}$ | 401 |
| If $0 < (\nu_1)^{(3)} < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\bar{\nu}_1)^{(3)}$ we find like in the previous case, | 402 |

$$(\nu_1)^{(3)} \leq \frac{(\nu_1)^{(3)}+(C)^{(3)}(\nu_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}}{1+(C)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}} \leq \nu^{(3)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\bar{\nu}_1)^{(3)}$$

If $0 < (\nu_1)^{(3)} \leq (\bar{\nu}_1)^{(3)} \leq (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0}$ , we obtain

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$$(\nu_1)^{(3)} \leq \nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\nu_0)^{(3)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(3)}(t)$ :-

$$(m_2)^{(3)} \leq \nu^{(3)}(t) \leq (m_1)^{(3)}, \quad \boxed{\nu^{(3)}(t) = \frac{G_{20}(t)}{G_{21}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(3)}(t)$ :-

$$(\mu_2)^{(3)} \leq u^{(3)}(t) \leq (\mu_1)^{(3)}, \quad \boxed{u^{(3)}(t) = \frac{T_{20}(t)}{T_{21}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{20}^{''})^{(3)} = (a_{21}^{''})^{(3)}, then\ (\sigma_1)^{(3)} = (\sigma_2)^{(3)}$ and in this case $(\nu_1)^{(3)} = (\bar{\nu}_1)^{(3)}$ if in addition $(\nu_0)^{(3)} = (\nu_1)^{(3)}$ then $\nu^{(3)}(t) = (\nu_0)^{(3)}$ and as a consequence $G_{20}(t) = (\nu_0)^{(3)}G_{21}(t)$

Analogously if $(b_{20}^{''})^{(3)} = (b_{21}^{''})^{(3)}, then\ (\tau_1)^{(3)} = (\tau_2)^{(3)}$ and then

$(u_1)^{(3)} = (\bar{u}_1)^{(3)}$if in addition $(u_0)^{(3)} = (u_1)^{(3)}$ then $T_{20}(t) = (u_0)^{(3)}T_{21}(t)$ This is an important consequence of the relation between $(\nu_1)^{(3)}$ and $(\bar{\nu}_1)^{(3)}$

**Proof** : From global equations we obtain

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$$\frac{d\nu^{(4)}}{dt} = (a_{24})^{(4)} - \left((a_{24}^{'})^{(4)} - (a_{25}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25},t)\right) - (a_{25}^{''})^{(4)}(T_{25},t)\nu^{(4)} - (a_{25})^{(4)}\nu^{(4)}$$

**Definition of** $\nu^{(4)}$ :- $\quad \boxed{\nu^{(4)} = \frac{G_{24}}{G_{25}}}$

It follows

$$-\left((a_{25})^{(4)}\left(\nu^{(4)}\right)^2 + (\sigma_2)^{(4)}\nu^{(4)} - (a_{24})^{(4)}\right) \leq \frac{d\nu^{(4)}}{dt} \leq -\left((a_{25})^{(4)}\left(\nu^{(4)}\right)^2 + (\sigma_4)^{(4)}\nu^{(4)} - (a_{24})^{(4)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(4)}, (\nu_0)^{(4)}$ :-

For $0 < \boxed{(\nu_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0}} < (\nu_1)^{(4)} < (\bar{\nu}_1)^{(4)}$

$$\nu^{(4)}(t) \geq \frac{(\nu_1)^{(4)} + (C)^{(4)}(\nu_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_0)^{(4)}\right)t\right]}}{4 + (C)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_0)^{(4)}\right)t\right]}} \quad , \quad \boxed{(C)^{(4)} = \frac{(\nu_1)^{(4)} - (\nu_0)^{(4)}}{(\nu_0)^{(4)} - (\nu_2)^{(4)}}}$$

it follows $(\nu_0)^{(4)} \leq \nu^{(4)}(t) \leq (\nu_1)^{(4)}$

In the same manner , we get

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$$\nu^{(4)}(t) \leq \frac{(\bar{\nu}_1)^{(4)} + (\bar{C})^{(4)}(\bar{\nu}_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}}{4 + (\bar{C})^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(4)} = \frac{(\bar{\nu}_1)^{(4)} - (\nu_0)^{(4)}}{(\nu_0)^{(4)} - (\bar{\nu}_2)^{(4)}}}$$

From which we deduce $(\nu_0)^{(4)} \leq \nu^{(4)}(t) \leq (\bar{\nu}_1)^{(4)}$

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If $0 < (\nu_1)^{(4)} < (\nu_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0} < (\bar{\nu}_1)^{(4)}$ we find like in the previous case,

$$(\nu_1)^{(4)} \leq \frac{(\nu_1)^{(4)} + (C)^{(4)}(\nu_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_2)^{(4)}\right)t\right]}}{1 + (C)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\nu_1)^{(4)} - (\nu_2)^{(4)}\right)t\right]}} \leq \nu^{(4)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(4)} + (\bar{C})^{(4)}(\bar{\nu}_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}}{1 + (\bar{C})^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}} \leq (\bar{\nu}_1)^{(4)}$$

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If $0 < (\nu_1)^{(4)} \leq (\bar{\nu}_1)^{(4)} \leq \boxed{(\nu_0)^{(4)} = \frac{G_{24}^0}{G_{25}^0}}$ , we obtain

$$(\nu_1)^{(4)} \leq \nu^{(4)}(t) \leq \frac{(\bar{\nu}_1)^{(4)} + (\bar{C})^{(4)}(\bar{\nu}_2)^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}}{1 + (\bar{C})^{(4)} e^{\left[-(a_{25})^{(4)}\left((\bar{\nu}_1)^{(4)} - (\bar{\nu}_2)^{(4)}\right)t\right]}} \leq (\nu_0)^{(4)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(4)}(t)$ :-

$(m_2)^{(4)} \leq \nu^{(4)}(t) \leq (m_1)^{(4)}$, $\boxed{\nu^{(4)}(t) = \frac{G_{24}(t)}{G_{25}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(4)}(t)$ :-

$(\mu_2)^{(4)} \leq u^{(4)}(t) \leq (\mu_1)^{(4)}$, $\boxed{u^{(4)}(t) = \frac{T_{24}(t)}{T_{25}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{24}^{''})^{(4)} = (a_{25}^{''})^{(4)}, then\ (\sigma_1)^{(4)} = (\sigma_2)^{(4)}$ and in this case $(\nu_1)^{(4)} = (\bar{\nu}_1)^{(4)}$ if in addition $(\nu_0)^{(4)} =$

$(\nu_1)^{(4)}$ then $\nu^{(4)}(t) = (\nu_0)^{(4)}$ and as a consequence $G_{24}(t) = (\nu_0)^{(4)} G_{25}(t)$ **this also defines** $(\nu_0)^{(4)}$ **for the special case** .

Analogously if $(b_{24}^{''})^{(4)} = (b_{25}^{''})^{(4)}, then\ (\tau_1)^{(4)} = (\tau_2)^{(4)}$ and then $(u_1)^{(4)} = (\bar{u}_4)^{(4)}$ if in addition $(u_0)^{(4)} = (u_1)^{(4)}$ then $T_{24}(t) = (u_0)^{(4)} T_{25}(t)$ This is an important consequence of the relation between $(\nu_1)^{(4)}$ and $(\bar{\nu}_1)^{(4)}$, **and definition of** $(u_0)^{(4)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(5)}}{dt} = (a_{28})^{(5)} - \left((a_{28}^{'})^{(5)} - (a_{29}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}, t)\right) - (a_{29}^{''})^{(5)}(T_{29}, t)\nu^{(5)} - (a_{29})^{(5)}\nu^{(5)}$$

**Definition of** $\nu^{(5)}$ :- $\boxed{\nu^{(5)} = \frac{G_{28}}{G_{29}}}$

It follows

$$-\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_2)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right) \leq \frac{d\nu^{(5)}}{dt} \leq -\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_1)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(5)}, (\nu_0)^{(5)}$ :-

For $0 < \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}} < (\nu_1)^{(5)} < (\bar{\nu}_1)^{(5)}$

$$\nu^{(5)}(t) \geq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}}{5 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}} \quad , \quad \boxed{(C)^{(5)} = \frac{(\nu_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\nu_2)^{(5)}}}$$

it follows $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\nu_1)^{(5)}$

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In the same manner , we get

$$\nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)} + (\bar{C})^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{5 + (\bar{C})^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(5)} = \frac{(\bar{\nu}_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\bar{\nu}_2)^{(5)}}}$$

From which we deduce $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\bar{\nu}_5)^{(5)}$

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If $0 < (\nu_1)^{(5)} < (\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0} < (\bar{\nu}_1)^{(5)}$ we find like in the previous case,

$$(\nu_1)^{(5)} \leq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}}{1 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}} \leq \nu^{(5)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(5)} + (\bar{C})^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{1 + (\bar{C})^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\bar{\nu}_1)^{(5)}$$

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If $0 < (\nu_1)^{(5)} \leq (\bar{\nu}_1)^{(5)} \leq \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}}$ , we obtain

$$(\nu_1)^{(5)} \leq \nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)}+(\bar{C})^{(5)}(\bar{\nu}_2)^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}}{1+(\bar{C})^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\nu_0)^{(5)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(5)}(t)$ :-

$$(m_2)^{(5)} \leq \nu^{(5)}(t) \leq (m_1)^{(5)}, \quad \boxed{\nu^{(5)}(t) = \frac{G_{28}(t)}{G_{29}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(5)}(t)$ :-

$$(\mu_2)^{(5)} \leq u^{(5)}(t) \leq (\mu_1)^{(5)}, \quad \boxed{u^{(5)}(t) = \frac{T_{28}(t)}{T_{29}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{28}^{''})^{(5)} = (a_{29}^{''})^{(5)}, then\ (\sigma_1)^{(5)} = (\sigma_2)^{(5)}$ and in this case $(\nu_1)^{(5)} = (\bar{\nu}_1)^{(5)}$ if in addition $(\nu_0)^{(5)} = (\nu_5)^{(5)}$ then $\nu^{(5)}(t) = (\nu_0)^{(5)}$ and as a consequence $G_{28}(t) = (\nu_0)^{(5)}G_{29}(t)$ **this also defines** $(\nu_0)^{(5)}$ **for the special case** .

Analogously if $(b_{28}^{''})^{(5)} = (b_{29}^{''})^{(5)}, then\ (\tau_1)^{(5)} = (\tau_2)^{(5)}$ and then $(u_1)^{(5)} = (\bar{u}_1)^{(5)}$if in addition $(u_0)^{(5)} = (u_1)^{(5)}$ then $T_{28}(t) = (u_0)^{(5)}T_{29}(t)$ This is an important consequence of the relation between $(\nu_1)^{(5)}$ and $(\bar{\nu}_1)^{(5)}$, **and definition of** $(u_0)^{(5)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(6)}}{dt} = (a_{32})^{(6)} - \left((a_{32}^{'})^{(6)} - (a_{33}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33}, t)\right) - (a_{33}^{''})^{(6)}(T_{33}, t)\nu^{(6)} - (a_{33})^{(6)}\nu^{(6)}$$

**Definition of** $\nu^{(6)}$ :- $\quad \boxed{\nu^{(6)} = \frac{G_{32}}{G_{33}}}$

It follows

$$-\left((a_{33})^{(6)}\left(\nu^{(6)}\right)^2 + (\sigma_2)^{(6)}\nu^{(6)} - (a_{32})^{(6)}\right) \leq \frac{d\nu^{(6)}}{dt} \leq -\left((a_{33})^{(6)}\left(\nu^{(6)}\right)^2 + (\sigma_1)^{(6)}\nu^{(6)} - (a_{32})^{(6)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(6)}, (\nu_0)^{(6)}$ :-

For $0 < \boxed{(\nu_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0}} < (\nu_1)^{(6)} < (\bar{\nu}_1)^{(6)}$

$$\nu^{(6)}(t) \geq \frac{(\nu_1)^{(6)}+(C)^{(6)}(\nu_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_0)^{(6)}\right)t\right]}}{1+(C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_0)^{(6)}\right)t\right]}} \quad , \quad \boxed{(C)^{(6)} = \frac{(\nu_1)^{(6)}-(\nu_0)^{(6)}}{(\nu_0)^{(6)}-(\nu_2)^{(6)}}}$$

it follows $(\nu_0)^{(6)} \leq \nu^{(6)}(t) \leq (\nu_1)^{(6)}$

In the same manner , we get 413

$$\nu^{(6)}(t) \le \frac{(\bar{\nu}_1)^{(6)}+(\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}}{1+(\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(6)} = \frac{(\bar{\nu}_1)^{(6)}-(\nu_0)^{(6)}}{(\nu_0)^{(6)}-(\bar{\nu}_2)^{(6)}}}$$

From which we deduce $(\nu_0)^{(6)} \le \nu^{(6)}(t) \le (\bar{\nu}_1)^{(6)}$

414

If $0 < (\nu_1)^{(6)} < (\nu_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0} < (\bar{\nu}_1)^{(6)}$ we find like in the previous case,

$$(\nu_1)^{(6)} \le \frac{(\nu_1)^{(6)}+(C)^{(6)}(\nu_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_2)^{(6)}\right)t\right]}}{1+(C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_2)^{(6)}\right)t\right]}} \le \nu^{(6)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(6)}+(\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}}{1+(\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}} \le (\bar{\nu}_1)^{(6)}$$

415

If $0 < (\nu_1)^{(6)} \le (\bar{\nu}_1)^{(6)} \le \boxed{(\nu_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0}}$ , we obtain

$$(\nu_1)^{(6)} \le \nu^{(6)}(t) \le \frac{(\bar{\nu}_1)^{(6)}+(\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}}{1+(\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}} \le (\nu_0)^{(6)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(6)}(t)$ :-

$$(m_2)^{(6)} \le \nu^{(6)}(t) \le (m_1)^{(6)}, \quad \boxed{\nu^{(6)}(t) = \frac{G_{32}(t)}{G_{33}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(6)}(t)$ :-

$$(\mu_2)^{(6)} \le u^{(6)}(t) \le (\mu_1)^{(6)}, \quad \boxed{u^{(6)}(t) = \frac{T_{32}(t)}{T_{33}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{32}^{''})^{(6)} = (a_{33}^{''})^{(6)}, then\ (\sigma_1)^{(6)} = (\sigma_2)^{(6)}$ and in this case $(\nu_1)^{(6)} = (\bar{\nu}_1)^{(6)}$ if in addition $(\nu_0)^{(6)} = (\nu_1)^{(6)}$ then $\nu^{(6)}(t) = (\nu_0)^{(6)}$ and as a consequence $G_{32}(t) = (\nu_0)^{(6)}G_{33}(t)$ **this also defines** $(\nu_0)^{(6)}$ **for the special case** .

Analogously if $(b_{32}^{''})^{(6)} = (b_{33}^{''})^{(6)}, then\ (\tau_1)^{(6)} = (\tau_2)^{(6)}$ and then $(u_1)^{(6)} = (\bar{u}_1)^{(6)}$ if in addition $(u_0)^{(6)} = (u_1)^{(6)}$ then $T_{32}(t) = (u_0)^{(6)}T_{33}(t)$ This is an important consequence of the relation between $(\nu_1)^{(6)}$ and $(\bar{\nu}_1)^{(6)}$, **and definition of** $(u_0)^{(6)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(7)}}{dt} = (a_{36})^{(7)} - \left((a_{36}^{'})^{(7)} - (a_{37}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37},t)\right) - (a_{37}^{''})^{(7)}(T_{37},t)\nu^{(7)} - (a_{37})^{(7)}\nu^{(7)}$$

**Definition of** $\nu^{(7)}$ :- $\boxed{\nu^{(7)} = \frac{G_{36}}{G_{37}}}$

It follows

$$-\left((a_{37})^{(7)}\left(\nu^{(7)}\right)^2 + (\sigma_2)^{(7)}\nu^{(7)} - (a_{36})^{(7)}\right) \leq \frac{d\nu^{(7)}}{dt} \leq -\left((a_{37})^{(7)}\left(\nu^{(7)}\right)^2 + (\sigma_1)^{(7)}\nu^{(7)} - (a_{36})^{(7)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(7)}, (\nu_0)^{(7)}$ :-

For $0 < \boxed{(\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0}} < (\nu_1)^{(7)} < (\bar{\nu}_1)^{(7)}$

$$\nu^{(7)}(t) \geq \frac{(\nu_1)^{(7)} + (C)^{(7)}(\nu_2)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)} - (\nu_0)^{(7)}\right)t\right]}}{1 + (C)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)} - (\nu_0)^{(7)}\right)t\right]}} \quad , \quad \boxed{(C)^{(7)} = \frac{(\nu_1)^{(7)} - (\nu_0)^{(7)}}{(\nu_0)^{(7)} - (\nu_2)^{(7)}}}$$

it follows $(\nu_0)^{(7)} \leq \nu^{(7)}(t) \leq (\nu_1)^{(7)}$

In the same manner , we get 417

$$\nu^{(7)}(t) \leq \frac{(\bar{\nu}_1)^{(7)} + (\bar{C})^{(7)}(\bar{\nu}_2)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)} - (\bar{\nu}_2)^{(7)}\right)t\right]}}{1 + (\bar{C})^{(7)} e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)} - (\bar{\nu}_2)^{(7)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(7)} = \frac{(\bar{\nu}_1)^{(7)} - (\nu_0)^{(7)}}{(\nu_0)^{(7)} - (\bar{\nu}_2)^{(7)}}}$$

From which we deduce $(\nu_0)^{(7)} \leq \nu^{(7)}(t) \leq (\bar{\nu}_1)^{(7)}$

If $0 < (\nu_1)^{(7)} < (\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0} < (\bar{\nu}_1)^{(7)}$ we find like in the previous case, 418

$$(\nu_1)^{(7)} \leq \frac{(\nu_1)^{(7)} + (C)^{(7)}(\nu_2)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)} - (\nu_2)^{(7)}\right)t\right]}}{1 + (C)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\nu_1)^{(7)} - (\nu_2)^{(7)}\right)t\right]}} \leq \nu^{(7)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(7)} + (\bar{C})^{(7)}(\bar{\nu}_2)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)} - (\bar{\nu}_2)^{(7)}\right)t\right]}}{1 + (\bar{C})^{(7)} e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)} - (\bar{\nu}_2)^{(7)}\right)t\right]}} \leq (\bar{\nu}_1)^{(7)}$$

If $0 < (\nu_1)^{(7)} \leq (\bar{\nu}_1)^{(7)} \leq \boxed{(\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0}}$ , we obtain 419

$$(\nu_1)^{(7)} \leq \nu^{(7)}(t) \leq \frac{(\bar{\nu}_1)^{(7)} + (\bar{C})^{(7)}(\bar{\nu}_2)^{(7)} e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)} - (\bar{\nu}_2)^{(7)}\right)t\right]}}{1 + (\bar{C})^{(7)} e^{\left[-(a_{37})^{(7)}\left((\bar{\nu}_1)^{(7)} - (\bar{\nu}_2)^{(7)}\right)t\right]}} \leq (\nu_0)^{(7)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(7)}(t)$ :-

$$(m_2)^{(7)} \leq \nu^{(7)}(t) \leq (m_1)^{(7)}, \quad \boxed{\nu^{(7)}(t) = \frac{G_{36}(t)}{G_{37}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(7)}(t)$ :- 420

$(\mu_2)^{(7)} \leq u^{(7)}(t) \leq (\mu_1)^{(7)}$, $\boxed{u^{(7)}(t) = \frac{T_{36}(t)}{T_{37}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{36}^{''})^{(7)} = (a_{37}^{''})^{(7)}, then\ (\sigma_1)^{(7)} = (\sigma_2)^{(7)}$ and in this case $(\nu_1)^{(7)} = (\bar{\nu}_1)^{(7)}$ if in addition $(\nu_0)^{(7)} = (\nu_1)^{(7)}$ then $\nu^{(7)}(t) = (\nu_0)^{(7)}$ and as a consequence $G_{36}(t) = (\nu_0)^{(7)} G_{37}(t)$ **this also defines** $(\nu_0)^{(7)}$ **for the special case** .

Analogously if $(b_{36}^{''})^{(7)} = (b_{37}^{''})^{(7)}, then\ (\tau_1)^{(7)} = (\tau_2)^{(7)}$ and then $(u_1)^{(7)} = (\bar{u}_1)^{(7)}$if in addition $(u_0)^{(7)} = (u_1)^{(7)}$ then $T_{36}(t) = (u_0)^{(7)} T_{37}(t)$ This is an important consequence of the relation between $(\nu_1)^{(7)}$ and $(\bar{\nu}_1)^{(7)}$, **and definition of** $(u_0)^{(7)}$.

**Proof** : From global equations we obtain 421

$$\frac{d\nu^{(8)}}{dt} = (a_{40})^{(8)} - \left((a_{40}^{'})^{(8)} - (a_{41}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41}, t)\right) - (a_{41}^{''})^{(8)}(T_{41}, t)\nu^{(8)} - (a_{41})^{(8)}\nu^{(8)}$$

**Definition of** $\nu^{(8)}$ :- $\boxed{\nu^{(8)} = \frac{G_{40}}{G_{41}}}$

It follows

$$-\left((a_{41})^{(8)}\left(\nu^{(8)}\right)^2 + (\sigma_2)^{(8)}\nu^{(8)} - (a_{40})^{(8)}\right) \leq \frac{d\nu^{(8)}}{dt} \leq -\left((a_{41})^{(8)}\left(\nu^{(8)}\right)^2 + (\sigma_1)^{(8)}\nu^{(8)} - (a_{40})^{(8)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(8)}, (\nu_0)^{(8)}$ :-

For $0 < \boxed{(\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}} < (\nu_1)^{(8)} < (\bar{\nu}_1)^{(8)}$

$$\nu^{(8)}(t) \geq \frac{(\nu_1)^{(8)} + (C)^{(8)}(\nu_2)^{(8)} e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}}{1 + (C)^{(8)} e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}} \quad , \quad \boxed{(C)^{(8)} = \frac{(\nu_1)^{(8)} - (\nu_0)^{(8)}}{(\nu_0)^{(8)} - (\nu_2)^{(8)}}}$$

it follows $(\nu_0)^{(8)} \leq \nu^{(8)}(t) \leq (\nu_1)^{(8)}$

In the same manner , we get 422

$$\nu^{(8)}(t) \le \frac{(\bar{\nu}_1)^{(8)}+(\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}}{1+(\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(8)} = \frac{(\bar{\nu}_1)^{(8)}-(\nu_0)^{(8)}}{(\nu_0)^{(8)}-(\bar{\nu}_2)^{(8)}}}$$

From which we deduce $(\nu_0)^{(8)} \le \nu^{(8)}(t) \le (\bar{\nu}_8)^{(8)}$

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If $0 < (\nu_1)^{(8)} < (\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0} < (\bar{\nu}_1)^{(8)}$ we find like in the previous case,

$$(\nu_1)^{(8)} \le \frac{(\nu_1)^{(8)}+(C)^{(8)}(\nu_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)}-(\nu_2)^{(8)}\right)t\right]}}{1+(C)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)}-(\nu_2)^{(8)}\right)t\right]}} \le \nu^{(8)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(8)}+(\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}}{1+(\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}} \le (\bar{\nu}_1)^{(8)}$$

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If $0 < (\nu_1)^{(8)} \le (\bar{\nu}_1)^{(8)} \le \boxed{(\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}}$, we obtain

$$(\nu_1)^{(8)} \le \nu^{(8)}(t) \le \frac{(\bar{\nu}_1)^{(8)}+(\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}}{1+(\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}} \le (\nu_0)^{(8)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(8)}(t)$ :-

$$(m_2)^{(8)} \le \nu^{(8)}(t) \le (m_1)^{(8)}, \quad \boxed{\nu^{(8)}(t) = \frac{G_{40}(t)}{G_{41}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(8)}(t)$ :-

$$(\mu_2)^{(8)} \le u^{(8)}(t) \le (\mu_1)^{(8)}, \quad \boxed{u^{(8)}(t) = \frac{T_{40}(t)}{T_{41}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{40}^{''})^{(8)} = (a_{41}^{''})^{(8)}, then\ (\sigma_1)^{(8)} = (\sigma_2)^{(8)}$ and in this case $(\nu_1)^{(8)} = (\bar{\nu}_1)^{(8)}$ if in addition $(\nu_0)^{(8)} = (\nu_1)^{(8)}$ then $\nu^{(8)}(t) = (\nu_0)^{(8)}$ and as a consequence $G_{40}(t) = (\nu_0)^{(8)}G_{41}(t)$ **this also defines** $(\nu_0)^{(8)}$ **for the special case** .

Analogously if $(b_{40}^{''})^{(8)} = (b_{41}^{''})^{(8)}, then\ (\tau_1)^{(8)} = (\tau_2)^{(8)}$ and then $(u_1)^{(8)} = (\bar{u}_1)^{(8)}$ if in addition $(u_0)^{(8)} = (u_1)^{(8)}$ then $T_{40}(t) = (u_0)^{(8)}T_{41}(t)$ This is an important consequence of the relation between $(\nu_1)^{(8)}$ and $(\bar{\nu}_1)^{(8)}$, **and definition of** $(u_0)^{(8)}$.

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**Proof** : From 99,20,44,22,23,44 we obtain

$$\frac{dv^{(9)}}{dt} = (a_{44})^{(9)} - \left((a'_{44})^{(9)} - (a'_{45})^{(9)} + (a''_{44})^{(9)}(T_{45},t)\right) - (a''_{45})^{(9)}(T_{45},t)v^{(9)} - (a_{45})^{(9)}v^{(9)}$$

**Definition of** $v^{(9)}$ :- $\boxed{v^{(9)} = \frac{G_{44}}{G_{45}}}$

It follows

$$-\left((a_{45})^{(9)}\left(v^{(9)}\right)^2 + (\sigma_2)^{(9)}v^{(9)} - (a_{44})^{(9)}\right) \le \frac{dv^{(9)}}{dt} \le -\left((a_{45})^{(9)}\left(v^{(9)}\right)^2 + (\sigma_1)^{(9)}v^{(9)} - (a_{44})^{(9)}\right)$$

From which one obtains

**Definition of** $(\bar{v}_1)^{(9)}, (v_0)^{(9)}$ :-

For $0 < \boxed{(v_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}}} < (v_1)^{(9)} < (\bar{v}_1)^{(9)}$

$$v^{(9)}(t) \ge \frac{(v_1)^{(9)} + (C)^{(9)}(v_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((v_1)^{(9)} - (v_0)^{(9)}\right)t\right]}}{1 + (C)^{(9)}e^{\left[-(a_{45})^{(9)}\left((v_1)^{(9)} - (v_0)^{(9)}\right)t\right]}} \quad , \quad \boxed{(C)^{(9)} = \frac{(v_1)^{(9)} - (v_0)^{(9)}}{(v_0)^{(9)} - (v_2)^{(9)}}}$$

it follows $(v_0)^{(9)} \le v^{(9)}(t) \le (v_9)^{(9)}$

In the same manner , we get

$$v^{(9)}(t) \le \frac{(\bar{v}_1)^{(9)} + (\bar{C})^{(9)}(\bar{v}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)} - (\bar{v}_2)^{(9)}\right)t\right]}}{1 + (\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)} - (\bar{v}_2)^{(9)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(9)} = \frac{(\bar{v}_1)^{(9)} - (v_0)^{(9)}}{(v_0)^{(9)} - (\bar{v}_2)^{(9)}}}$$

From which we deduce $(v_0)^{(9)} \le v^{(9)}(t) \le (\bar{v}_1)^{(9)}$

If $0 < (v_1)^{(9)} < (v_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}} < (\bar{v}_1)^{(9)}$ we find like in the previous case,

$$(v_1)^{(9)} \le \frac{(v_1)^{(9)} + (C)^{(9)}(v_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((v_1)^{(9)} - (v_2)^{(9)}\right)t\right]}}{1 + (C)^{(9)}e^{\left[-(a_{45})^{(9)}\left((v_1)^{(9)} - (v_2)^{(9)}\right)t\right]}} \le v^{(9)}(t) \le$$

$$\frac{(\bar{v}_1)^{(9)} + (\bar{C})^{(9)}(\bar{v}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)} - (\bar{v}_2)^{(9)}\right)t\right]}}{1 + (\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)} - (\bar{v}_2)^{(9)}\right)t\right]}} \le (\bar{v}_1)^{(9)}$$

If $0 < (v_1)^{(9)} \le (\bar{v}_1)^{(9)} \le \boxed{(v_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}}}$ , we obtain

$$(v_1)^{(9)} \le v^{(9)}(t) \le \frac{(\bar{v}_1)^{(9)} + (\bar{C})^{(9)}(\bar{v}_2)^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)} - (\bar{v}_2)^{(9)}\right)t\right]}}{1 + (\bar{C})^{(9)}e^{\left[-(a_{45})^{(9)}\left((\bar{v}_1)^{(9)} - (\bar{v}_2)^{(9)}\right)t\right]}} \le (v_0)^{(9)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $v^{(9)}(t)$ :-

$$(m_2)^{(9)} \le v^{(9)}(t) \le (m_1)^{(9)}, \quad \boxed{v^{(9)}(t) = \frac{G_{44}(t)}{G_{45}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(9)}(t)$ :-

$(\mu_2)^{(9)} \leq u^{(9)}(t) \leq (\mu_1)^{(9)}$, $\boxed{u^{(9)}(t) = \frac{T_{44}(t)}{T_{45}(t)}}$

Now, using this result and replacing it in 99, 20,44,22,23, and 44 we get easily the result stated in the theorem.

**Particular case :**

If $(a_{44}^{''})^{(9)} = (a_{45}^{''})^{(9)}, then\ (\sigma_1)^{(9)} = (\sigma_2)^{(9)}$ and in this case $(\nu_1)^{(9)} = (\bar{\nu}_1)^{(9)}$ if in addition $(\nu_0)^{(9)} = (\nu_1)^{(9)}$ then $\nu^{(9)}(t) = (\nu_0)^{(9)}$ and as a consequence $G_{44}(t) = (\nu_0)^{(9)} G_{45}(t)$ **this also defines** $(\nu_0)^{(9)}$ **for the special case** .

Analogously if $(b_{44}^{''})^{(9)} = (b_{45}^{''})^{(9)}, then\ (\tau_1)^{(9)} = (\tau_2)^{(9)}$ and then $(u_1)^{(9)} = (\bar{u}_1)^{(9)}$if in addition $(u_0)^{(9)} = (u_1)^{(9)}$ then $T_{44}(t) = (u_0)^{(9)} T_{45}(t)$ This is an important consequence of the relation between $(\nu_1)^{(9)}$ and $(\bar{\nu}_1)^{(9)}$, **and definition of** $(u_0)^{(9)}$.

We can prove the following 425

**Theorem :** If $(a_i^{''})^{(1)} and\ (b_i^{''})^{(1)}$ are independent on $t$ , and the conditions with the notations

$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} < 0$

$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a_{13})^{(1)}(p_{13})^{(1)} + (a_{14}^{'})^{(1)}(p_{14})^{(1)} + (p_{13})^{(1)}(p_{14})^{(1)} > 0$

$(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} > 0$ ,

$(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} - (b_{13}^{'})^{(1)}(r_{14})^{(1)} - (b_{14}^{'})^{(1)}(r_{14})^{(1)} + (r_{13})^{(1)}(r_{14})^{(1)} < 0$

$with\ (p_{13})^{(1)}, (r_{14})^{(1)}$ as defined by equation are satisfied , then the system

**Theorem :** If $(a_i^{''})^{(2)} and\ (b_i^{''})^{(2)}$ are independent on t , and the conditions with the notations 426

$(a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} < 0$ 427

$(a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a_{16})^{(2)}(p_{16})^{(2)} + (a_{17}^{'})^{(2)}(p_{17})^{(2)} + (p_{16})^{(2)}(p_{17})^{(2)} > 0$ 428

$(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} > 0$ , 429

$(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} - (b_{16}^{'})^{(2)}(r_{17})^{(2)} - (b_{17}^{'})^{(2)}(r_{17})^{(2)} + (r_{16})^{(2)}(r_{17})^{(2)} < 0$ 430

$with\ (p_{16})^{(2)}, (r_{17})^{(2)}$ as defined by equation are satisfied , then the system

**Theorem :** If $(a_i^{''})^{(3)} and\ (b_i^{''})^{(3)}$ are independent on $t$ , and the conditions with the notations 431

$(a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} < 0$

$(a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a_{20})^{(3)}(p_{20})^{(3)} + (a_{21}^{'})^{(3)}(p_{21})^{(3)} + (p_{20})^{(3)}(p_{21})^{(3)} > 0$

$(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} > 0$ ,

$(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} - (b_{20}^{'})^{(3)}(r_{21})^{(3)} - (b_{21}^{'})^{(3)}(r_{21})^{(3)} + (r_{20})^{(3)}(r_{21})^{(3)} < 0$

| | |
|---|---|
| $with\ \ (p_{20})^{(3)}, (r_{21})^{(3)}$ as defined by equation are satisfied , then the system | |
| We can prove the following<br><br>**Theorem :** If $(a_i^{''})^{(4)} and\ (b_i^{''})^{(4)}$ are independent on $t$ , and the conditions with the notations<br><br>$(a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} < 0$<br><br>$(a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a_{24})^{(4)}(p_{24})^{(4)} + (a_{25}^{'})^{(4)}(p_{25})^{(4)} + (p_{24})^{(4)}(p_{25})^{(4)} > 0$<br><br>$(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} > 0$ ,<br><br>$(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} - (b_{24}^{'})^{(4)}(r_{25})^{(4)} - (b_{25}^{'})^{(4)}(r_{25})^{(4)} + (r_{24})^{(4)}(r_{25})^{(4)} < 0$<br><br>$with\ \ (p_{24})^{(4)}, (r_{25})^{(4)}$ as defined by equation are satisfied , then the system | 432 |
| **Theorem :** If $(a_i^{''})^{(5)} and\ (b_i^{''})^{(5)}$ are independent on $t$ , and the conditions with the notations<br><br>$(a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} < 0$<br><br>$(a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a_{28})^{(5)}(p_{28})^{(5)} + (a_{29}^{'})^{(5)}(p_{29})^{(5)} + (p_{28})^{(5)}(p_{29})^{(5)} > 0$<br><br>$(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} > 0$ ,<br><br>$(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} - (b_{28}^{'})^{(5)}(r_{29})^{(5)} - (b_{29}^{'})^{(5)}(r_{29})^{(5)} + (r_{28})^{(5)}(r_{29})^{(5)} < 0$<br><br>$with\ \ (p_{28})^{(5)}, (r_{29})^{(5)}$ as defined by equation are satisfied , then the system | 433 |
| **Theorem** If $(a_i^{''})^{(6)} and\ (b_i^{''})^{(6)}$ are independent on $t$ , and the conditions with the notations<br><br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} < 0$<br><br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32})^{(6)}(p_{32})^{(6)} + (a_{33}^{'})^{(6)}(p_{33})^{(6)} + (p_{32})^{(6)}(p_{33})^{(6)} > 0$<br><br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} > 0$ ,<br><br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} - (b_{32}^{'})^{(6)}(r_{33})^{(6)} - (b_{33}^{'})^{(6)}(r_{33})^{(6)} + (r_{32})^{(6)}(r_{33})^{(6)} < 0$<br><br>$with\ \ (p_{32})^{(6)}, (r_{33})^{(6)}$ as defined by equation are satisfied , then the system | 434 |
| **Theorem :** If $(a_i^{''})^{(7)} and\ (b_i^{''})^{(7)}$ are independent on $t$ , and the conditions with the notations<br><br>$(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} < 0$<br><br>$(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36})^{(7)}(p_{36})^{(7)} + (a_{37}^{'})^{(7)}(p_{37})^{(7)} + (p_{36})^{(7)}(p_{37})^{(7)} > 0$<br><br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} > 0$ ,<br><br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} - (b_{36}^{'})^{(7)}(r_{37})^{(7)} - (b_{37}^{'})^{(7)}(r_{37})^{(7)} + (r_{36})^{(7)}(r_{37})^{(7)} < 0$ | 435 |

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|---|---|
| $with\ \ (p_{36})^{(7)}, (r_{37})^{(7)}$ as defined by equation are satisfied , then the system | |
| **Theorem :** If $(a_i^{''})^{(8)} and\ (b_i^{''})^{(8)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} < 0$<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40})^{(8)}(p_{40})^{(8)} + (a_{41}^{'})^{(8)}(p_{41})^{(8)} + (p_{40})^{(8)}(p_{41})^{(8)} > 0$<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} > 0$ ,<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} - (b_{40}^{'})^{(8)}(r_{41})^{(8)} - (b_{41}^{'})^{(8)}(r_{41})^{(8)} + (r_{40})^{(8)}(r_{41})^{(8)} < 0$<br>$with\ \ (p_{40})^{(8)}, (r_{41})^{(8)}$ as defined by equation are satisfied , then the system | 436 |
| **Theorem :** If $(a_i^{''})^{(9)} and\ (b_i^{''})^{(9)}$ are independent on $t$ , and the conditions (with the notations 45,46,27,28)<br>$(a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} < 0$<br>$(a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44})^{(9)}(p_{44})^{(9)} + (a_{45}^{'})^{(9)}(p_{45})^{(9)} + (p_{44})^{(9)}(p_{45})^{(9)} > 0$<br>$(b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} > 0$ ,<br>$(b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} - (b_{44}^{'})^{(9)}(r_{45})^{(9)} - (b_{45}^{'})^{(9)}(r_{45})^{(9)} + (r_{44})^{(9)}(r_{45})^{(9)} < 0$<br>$with\ \ (p_{44})^{(9)}, (r_{45})^{(9)}$ as defined by equation 45 are satisfied , then the system | 436 A |
| $(a_{13})^{(1)}G_{14} - \left[(a_{13}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14})\right]G_{13} = 0$ | 437 |
| $(a_{14})^{(1)}G_{13} - \left[(a_{14}^{'})^{(1)} + (a_{14}^{''})^{(1)}(T_{14})\right]G_{14} = 0$ | 438 |
| $(a_{15})^{(1)}G_{14} - \left[(a_{15}^{'})^{(1)} + (a_{15}^{''})^{(1)}(T_{14})\right]G_{15} = 0$ | 439 |
| $(b_{13})^{(1)}T_{14} - [(b_{13}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G)\ ]T_{13} = 0$ | 440 |
| $(b_{14})^{(1)}T_{13} - [(b_{14}^{'})^{(1)} - (b_{14}^{''})^{(1)}(G)\ ]T_{14} = 0$ | 441 |
| $(b_{15})^{(1)}T_{14} - [(b_{15}^{'})^{(1)} - (b_{15}^{''})^{(1)}(G)\ ]T_{15} = 0$ | 442 |
| has a unique positive solution , which is an equilibrium solution for the system | |
| $(a_{16})^{(2)}G_{17} - \left[(a_{16}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17})\right]G_{16} = 0$ | 443 |
| $(a_{17})^{(2)}G_{16} - \left[(a_{17}^{'})^{(2)} + (a_{17}^{''})^{(2)}(T_{17})\right]G_{17} = 0$ | 444 |
| $(a_{18})^{(2)}G_{17} - \left[(a_{18}^{'})^{(2)} + (a_{18}^{''})^{(2)}(T_{17})\right]G_{18} = 0$ | 445 |

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| $(b_{16})^{(2)}T_{17} - [(b_{16}^{'})^{(2)} - (b_{16}^{''})^{(2)}(G_{19})\,]T_{16} = 0$ | 446 |
| $(b_{17})^{(2)}T_{16} - [(b_{17}^{'})^{(2)} - (b_{17}^{''})^{(2)}(G_{19})\,]T_{17} = 0$ | 447 |
| $(b_{18})^{(2)}T_{17} - [(b_{18}^{'})^{(2)} - (b_{18}^{''})^{(2)}(G_{19})\,]T_{18} = 0$ | 448 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{20})^{(3)}G_{21} - \left[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21})\right]G_{20} = 0$ | 449 |
| $(a_{21})^{(3)}G_{20} - \left[(a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)}(T_{21})\right]G_{21} = 0$ | 450 |
| $(a_{22})^{(3)}G_{21} - \left[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21})\right]G_{22} = 0$ | 451 |
| $(b_{20})^{(3)}T_{21} - [(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23})\,]T_{20} = 0$ | 452 |
| $(b_{21})^{(3)}T_{20} - [(b_{21}^{'})^{(3)} - (b_{21}^{''})^{(3)}(G_{23})\,]T_{21} = 0$ | 453 |
| $(b_{22})^{(3)}T_{21} - [(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23})\,]T_{22} = 0$ | 454 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25})\right]G_{24} = 0$ | 455 |
| $(a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25})\right]G_{25} = 0$ | 456 |
| $(a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25})\right]G_{26} = 0$ | 457 |
| $(b_{24})^{(4)}T_{25} - [(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}\big((G_{27})\big)\,]T_{24} = 0$ | 458 |
| $(b_{25})^{(4)}T_{24} - [(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}\big((G_{27})\big)\,]T_{25} = 0$ | 459 |
| $(b_{26})^{(4)}T_{25} - [(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}\big((G_{27})\big)\,]T_{26} = 0$ | 460 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29})\right]G_{28} = 0$ | 461 |
| $(a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29})\right]G_{29} = 0$ | 462 |
| $(a_{30})^{(5)}G_{29} - \left[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29})\right]G_{30} = 0$ | 463 |
| $(b_{28})^{(5)}T_{29} - [(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}(G_{31})\,]T_{28} = 0$ | 464 |
| $(b_{29})^{(5)}T_{28} - [(b_{29}^{'})^{(5)} - (b_{29}^{''})^{(5)}(G_{31})\,]T_{29} = 0$ | 465 |

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| $(b_{30})^{(5)}T_{29} - [(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}(G_{31})\,]T_{30} = 0$ | 466 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{32})^{(6)}G_{33} - \left[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33})\right]G_{32} = 0$ | 467 |
| $(a_{33})^{(6)}G_{32} - \left[(a_{33}^{'})^{(6)} + (a_{33}^{''})^{(6)}(T_{33})\right]G_{33} = 0$ | 468 |
| $(a_{34})^{(6)}G_{33} - \left[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33})\right]G_{34} = 0$ | 469 |
| $(b_{32})^{(6)}T_{33} - [(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}(G_{35})\,]T_{32} = 0$ | 470 |
| $(b_{33})^{(6)}T_{32} - [(b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}(G_{35})\,]T_{33} = 0$ | 471 |
| $(b_{34})^{(6)}T_{33} - [(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}(G_{35})\,]T_{34} = 0$ | 472 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{36})^{(7)}G_{37} - \left[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37})\right]G_{36} = 0$ | 473 |
| $(a_{37})^{(7)}G_{36} - \left[(a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37})\right]G_{37} = 0$ | 474 |
| $(a_{38})^{(7)}G_{37} - \left[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37})\right]G_{38} = 0$ | 475 |
| $(b_{36})^{(7)}T_{37} - [(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}(G_{39})]T_{36} = 0$ | 476 |
| $(b_{37})^{(7)}T_{36} - [(b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}(G_{39})\,]T_{37} = 0$ | 477 |
| $(b_{38})^{(7)}T_{37} - [(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}(G_{39})\,]T_{38} = 0$ | 478 |
| $(a_{40})^{(8)}G_{41} - \left[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41})\right]G_{40} = 0$ | 479 |
| $(a_{41})^{(8)}G_{40} - \left[(a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41})\right]G_{41} = 0$ | 480 |
| $(a_{42})^{(8)}G_{41} - \left[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41})\right]G_{42} = 0$ | 481 |
| $(b_{40})^{(8)}T_{41} - [(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}(G_{43})\,]T_{40} = 0$ | 482 |
| $(b_{41})^{(8)}T_{40} - [(b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}(G_{43})\,]T_{41} = 0$ | 483 |

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| $(b_{42})^{(8)}T_{41} - [(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}(G_{43})\,]T_{42} = 0$ | 484 |
| $(a_{44})^{(9)}G_{45} - \left[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45})\right]G_{44} = 0$ | 484 A |
| $(a_{45})^{(9)}G_{44} - \left[(a_{45}^{'})^{(9)} + (a_{45}^{''})^{(9)}(T_{45})\right]G_{45} = 0$ | |
| $(a_{46})^{(9)}G_{45} - \left[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45})\right]G_{46} = 0$ | |
| $(b_{44})^{(9)}T_{45} - [(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}(G_{47})\,]T_{44} = 0$ | |
| $(b_{45})^{(9)}T_{44} - [(b_{45}^{'})^{(9)} - (b_{45}^{''})^{(9)}(G_{47})\,]T_{45} = 0$ | |
| $(b_{46})^{(9)}T_{45} - [(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}(G_{47})\,]T_{46} = 0$ | |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{13}, G_{14}$ if<br><br>$F(T) = (a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a_{13}^{'})^{(1)}(a_{14}^{''})^{(1)}(T_{14}) + (a_{14}^{'})^{(1)}(a_{13}^{''})^{(1)}(T_{14}) + (a_{13}^{''})^{(1)}(T_{14})(a_{14}^{''})^{(1)}(T_{14}) = 0$ | 485 |
| **Proof:**<br><br>(g) Indeed the first two equations have a nontrivial solution $G_{16}, G_{17}$ if<br><br>$\mathrm{F}(T_{19}) = (a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a_{16}^{'})^{(2)}(a_{17}^{''})^{(2)}(T_{17}) + (a_{17}^{'})^{(2)}(a_{16}^{''})^{(2)}(T_{17}) + (a_{16}^{''})^{(2)}(T_{17})(a_{17}^{''})^{(2)}(T_{17}) = 0$ | 486 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{20}, G_{21}$ if<br><br>$F(T_{23}) = (a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a_{20}^{'})^{(3)}(a_{21}^{''})^{(3)}(T_{21}) + (a_{21}^{'})^{(3)}(a_{20}^{''})^{(3)}(T_{21}) + (a_{20}^{''})^{(3)}(T_{21})(a_{21}^{''})^{(3)}(T_{21}) = 0$ | 487 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{24}, G_{25}$ if<br><br>$F(T_{27}) = (a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a_{24}^{'})^{(4)}(a_{25}^{''})^{(4)}(T_{25}) + (a_{25}^{'})^{(4)}(a_{24}^{''})^{(4)}(T_{25}) + (a_{24}^{''})^{(4)}(T_{25})(a_{25}^{''})^{(4)}(T_{25}) = 0$ | 488 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{28}, G_{29}$ if<br><br>$F(T_{31}) = (a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a_{28}^{'})^{(5)}(a_{29}^{''})^{(5)}(T_{29}) + (a_{29}^{'})^{(5)}(a_{28}^{''})^{(5)}(T_{29}) +$ | 489 |

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| $(a_{28}^{''})^{(5)}(T_{29})(a_{29}^{''})^{(5)}(T_{29}) = 0$ | |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{32}, G_{33}$ if<br><br>$F(T_{35}) = (a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32}^{'})^{(6)}(a_{33}^{''})^{(6)}(T_{33}) + (a_{33}^{'})^{(6)}(a_{32}^{''})^{(6)}(T_{33}) + (a_{32}^{''})^{(6)}(T_{33})(a_{33}^{''})^{(6)}(T_{33}) = 0$ | 490 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{36}, G_{37}$ if<br><br>$F(T_{39}) = (a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36}^{'})^{(7)}(a_{37}^{''})^{(7)}(T_{37}) + (a_{37}^{'})^{(7)}(a_{36}^{''})^{(7)}(T_{37}) + (a_{36}^{''})^{(7)}(T_{37})(a_{37}^{''})^{(7)}(T_{37}) = 0$ | 491 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{40}, G_{41}$ if<br><br>$F(T_{43}) = (a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40}^{'})^{(8)}(a_{41}^{''})^{(8)}(T_{41}) + (a_{41}^{'})^{(8)}(a_{40}^{''})^{(8)}(T_{41}) + (a_{40}^{''})^{(8)}(T_{41})(a_{41}^{''})^{(8)}(T_{41}) = 0$ | 492 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{44}, G_{45}$ if<br><br>$F(T_{47}) = (a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44}^{'})^{(9)}(a_{45}^{''})^{(9)}(T_{45}) + (a_{45}^{'})^{(9)}(a_{44}^{''})^{(9)}(T_{45}) + (a_{44}^{''})^{(9)}(T_{45})(a_{45}^{''})^{(9)}(T_{45}) = 0$ | 492 A |
| **Definition and uniqueness of** $T_{14}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{14})$ being increasing, it follows that there exists a unique $T_{14}^{*}$ for which $f(T_{14}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{13} = \frac{(a_{13})^{(1)}G_{14}}{[(a_{13}^{'})^{(1)}+(a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15} = \frac{(a_{15})^{(1)}G_{14}}{[(a_{15}^{'})^{(1)}+(a_{15}^{''})^{(1)}(T_{14}^{*})]}$ | 493 |
| **Definition and uniqueness of** $T_{17}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(2)}(T_{17})$ being increasing, it follows that there exists a unique $T_{17}^{*}$ for which $f(T_{17}^{*}) = 0$. With this value , we obtain from the three first equations | 494 |
| $G_{16} = \frac{(a_{16})^{(2)}G_{17}}{[(a_{16}^{'})^{(2)}+(a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18} = \frac{(a_{18})^{(2)}G_{17}}{[(a_{18}^{'})^{(2)}+(a_{18}^{''})^{(2)}(T_{17}^{*})]}$ | 495 |
| **Definition and uniqueness of** $T_{21}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{21})$ being increasing, it follows that | 496 |

there exists a unique $T_{21}^{*}$ for which $f(T_{21}^{*})=0$. With this value , we obtain from the three first equations

$$G_{20}=\frac{(a_{20})^{(3)}G_{21}}{[(a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21}^{*})]} \quad , \quad G_{22}=\frac{(a_{22})^{(3)}G_{21}}{[(a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21}^{*})]}$$

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**Definition and uniqueness of** $T_{25}^{*}$ :-

After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(4)}(T_{25})$ being increasing, it follows that there exists a unique $T_{25}^{*}$ for which $f(T_{25}^{*})=0$. With this value , we obtain from the three first equations

$$G_{24}=\frac{(a_{24})^{(4)}G_{25}}{[(a_{24}^{'})^{(4)}+(a_{24}^{''})^{(4)}(T_{25}^{*})]} \quad , \quad G_{26}=\frac{(a_{26})^{(4)}G_{25}}{[(a_{26}^{'})^{(4)}+(a_{26}^{''})^{(4)}(T_{25}^{*})]}$$

498

**Definition and uniqueness of** $T_{29}^{*}$ :-

After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(5)}(T_{29})$ being increasing, it follows that there exists a unique $T_{29}^{*}$ for which $f(T_{29}^{*})=0$. With this value , we obtain from the three first equations

$$G_{28}=\frac{(a_{28})^{(5)}G_{29}}{[(a_{28}^{'})^{(5)}+(a_{28}^{''})^{(5)}(T_{29}^{*})]} \quad , \quad G_{30}=\frac{(a_{30})^{(5)}G_{29}}{[(a_{30}^{'})^{(5)}+(a_{30}^{''})^{(5)}(T_{29}^{*})]}$$

499

**Definition and uniqueness of** $T_{33}^{*}$ :-

After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(6)}(T_{33})$ being increasing, it follows that there exists a unique $T_{33}^{*}$ for which $f(T_{33}^{*})=0$. With this value , we obtain from the three first equations

$$G_{32}=\frac{(a_{32})^{(6)}G_{33}}{[(a_{32}^{'})^{(6)}+(a_{32}^{''})^{(6)}(T_{33}^{*})]} \quad , \quad G_{34}=\frac{(a_{34})^{(6)}G_{33}}{[(a_{34}^{'})^{(6)}+(a_{34}^{''})^{(6)}(T_{33}^{*})]}$$

500

**Definition and uniqueness of** $T_{37}^{*}$ :-

After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(7)}(T_{37})$ being increasing, it follows that there exists a unique $T_{37}^{*}$ for which $f(T_{37}^{*})=0$. With this value , we obtain from the three first equations

$$G_{36}=\frac{(a_{36})^{(7)}G_{37}}{[(a_{36}^{'})^{(7)}+(a_{36}^{''})^{(7)}(T_{37}^{*})]} \quad , \quad G_{38}=\frac{(a_{38})^{(7)}G_{37}}{[(a_{38}^{'})^{(7)}+(a_{38}^{''})^{(7)}(T_{37}^{*})]}$$

501

**Definition and uniqueness of** $T_{41}^{*}$ :-

After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(8)}(T_{41})$ being increasing, it follows that there exists a unique $T_{41}^{*}$ for which $f(T_{41}^{*})=0$. With this value , we obtain from the three first equations

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| $G_{40} = \frac{(a_{40})^{(8)}G_{41}}{[(a_{40}^{'})^{(8)}+(a_{40}^{''})^{(8)}(T_{41}^{*})]}$ , $G_{42} = \frac{(a_{42})^{(8)}G_{41}}{[(a_{42}^{'})^{(8)}+(a_{42}^{''})^{(8)}(T_{41}^{*})]}$ | |
| **Definition and uniqueness of** $T_{45}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(9)}(T_{45})$ being increasing, it follows that there exists a unique $T_{45}^{*}$ for which $f(T_{45}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{44} = \frac{(a_{44})^{(9)}G_{45}}{[(a_{44}^{'})^{(9)}+(a_{44}^{''})^{(9)}(T_{45}^{*})]}$ , $G_{46} = \frac{(a_{46})^{(9)}G_{45}}{[(a_{46}^{'})^{(9)}+(a_{46}^{''})^{(9)}(T_{45}^{*})]}$ | 501 A |
| By the same argument, the equations admit solutions $G_{13}, G_{14}$ if<br>$\varphi(G) = (b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} -$<br>$\left[(b_{13}^{'})^{(1)}(b_{14}^{''})^{(1)}(G) + (b_{14}^{'})^{(1)}(b_{13}^{''})^{(1)}(G)\right]+(b_{13}^{''})^{(1)}(G)(b_{14}^{''})^{(1)}(G) = 0$<br><br>Where in $G(G_{13}, G_{14}, G_{15}), G_{13}, G_{15}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{14}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{14}^{*}$ such that $\varphi(G^{*}) = 0$ | 502 |
| By the same argument, the equations admit solutions $G_{16}, G_{17}$ if<br><br>$\varphi(G_{19}) = (b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} -$<br><br>$\left[(b_{16}^{'})^{(2)}(b_{17}^{''})^{(2)}(G_{19}) + (b_{17}^{'})^{(2)}(b_{16}^{''})^{(2)}(G_{19})\right]+(b_{16}^{''})^{(2)}(G_{19})(b_{17}^{''})^{(2)}(G_{19}) = 0$ | 503 |
| Where in $(G_{19})(G_{16}, G_{17}, G_{18}), G_{16}, G_{18}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{17}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{14}^{*}$ such that $\varphi((G_{19})^{*}) = 0$ | 504 |
| By the same argument, the equations admit solutions $G_{20}, G_{21}$ if<br><br>$\varphi(G_{23}) = (b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} -$<br><br>$\left[(b_{20}^{'})^{(3)}(b_{21}^{''})^{(3)}(G_{23}) + (b_{21}^{'})^{(3)}(b_{20}^{''})^{(3)}(G_{23})\right]+(b_{20}^{''})^{(3)}(G_{23})(b_{21}^{''})^{(3)}(G_{23}) = 0$<br><br>Where in $G_{23}(G_{20}, G_{21}, G_{22}), G_{20}, G_{22}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{21}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{21}^{*}$ such that $\varphi((G_{23})^{*}) = 0$ | 505 |
| By the same argument, the equations admit solutions $G_{24}, G_{25}$ if<br><br>$\varphi(G_{27}) = (b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} -$<br><br>$\left[(b_{24}^{'})^{(4)}(b_{25}^{''})^{(4)}(G_{27}) + (b_{25}^{'})^{(4)}(b_{24}^{''})^{(4)}(G_{27})\right]+(b_{24}^{''})^{(4)}(G_{27})(b_{25}^{''})^{(4)}(G_{27}) = 0$<br><br>Where in $(G_{27})(G_{24}, G_{25}, G_{26}), G_{24}, G_{26}$ must be replaced by their values from 96. It is easy to see that φ | 506 |

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| is a decreasing function in $G_{25}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{25}^{*}$ such that $\varphi((G_{27})^{*}) = 0$ | |
| By the same argument, the equations admit solutions $G_{28}, G_{29}$ if<br>$\varphi(G_{31}) = (b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} -$<br>$\left[(b_{28}^{'})^{(5)}(b_{29}^{''})^{(5)}(G_{31}) + (b_{29}^{'})^{(5)}(b_{28}^{''})^{(5)}(G_{31})\right] + (b_{28}^{''})^{(5)}(G_{31})(b_{29}^{''})^{(5)}(G_{31}) = 0$<br>Where in $(G_{31})(G_{28}, G_{29}, G_{30}), G_{28}, G_{30}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{29}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{29}^{*}$ such that $\varphi((G_{31})^{*}) = 0$ | 507 |
| By the same argument, the equations admit solutions $G_{32}, G_{33}$ if<br>$\varphi(G_{35}) = (b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} -$<br>$\left[(b_{32}^{'})^{(6)}(b_{33}^{''})^{(6)}(G_{35}) + (b_{33}^{'})^{(6)}(b_{32}^{''})^{(6)}(G_{35})\right] + (b_{32}^{''})^{(6)}(G_{35})(b_{33}^{''})^{(6)}(G_{35}) = 0$<br>Where in $(G_{35})(G_{32}, G_{33}, G_{34}), G_{32}, G_{34}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{33}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{33}^{*}$ such that $\varphi({G_{35}}^{*}) = 0$ | 508 |
| By the same argument, the equations admit solutions $G_{36}, G_{37}$ if<br>$\varphi(G_{39}) = (b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} -$<br>$\left[(b_{36}^{'})^{(7)}(b_{37}^{''})^{(7)}(G_{39}) + (b_{37}^{'})^{(7)}(b_{36}^{''})^{(7)}(G_{39})\right] + (b_{36}^{''})^{(7)}(G_{39})(b_{37}^{''})^{(7)}(G_{39}) = 0$<br>Where in $(G_{39})(G_{36}, G_{37}, G_{38}), G_{36}, G_{38}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{37}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{37}^{*}$ such that $\varphi({G_{39}}^{*}) = 0$ | 509 |
| By the same argument, the equations admit solutions $G_{40}, G_{41}$ if<br>$\varphi(G_{43}) = (b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} -$<br>$\left[(b_{40}^{'})^{(8)}(b_{41}^{''})^{(8)}(G_{43}) + (b_{41}^{'})^{(8)}(b_{40}^{''})^{(8)}(G_{43})\right] + (b_{40}^{''})^{(8)}(G_{43})(b_{41}^{''})^{(8)}(G_{43}) = 0$<br>Where in $(G_{43})(G_{40}, G_{41}, G_{42}), G_{40}, G_{42}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{41}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{41}^{*}$ such that $\varphi({G_{43}}^{*}) = 0$ | 510 |
| By the same argument, the equations 92,93 admit solutions $G_{44}, G_{45}$ if<br>$\varphi(G_{47}) = (b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} -$<br>$\left[(b_{44}^{'})^{(9)}(b_{45}^{''})^{(9)}(G_{47}) + (b_{45}^{'})^{(9)}(b_{44}^{''})^{(9)}(G_{47})\right] + (b_{44}^{''})^{(9)}(G_{47})(b_{45}^{''})^{(9)}(G_{47}) = 0$<br>Where in $(G_{47})(G_{44}, G_{45}, G_{46}), G_{44}, G_{46}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{45}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{45}^{*}$ such that $\varphi((G_{47})^{*}) = 0$ | |
| Finally we obtain the unique solution | 511 |

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| $G_{14}^{*}$ given by $\varphi(G^{*})=0$ , $T_{14}^{*}$ given by $f(T_{14}^{*})=0$ and<br>$G_{13}^{*}=\frac{(a_{13})^{(1)}G_{14}^{*}}{[(a_{13}^{'})^{(1)}+(a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15}^{*}=\frac{(a_{15})^{(1)}G_{14}^{*}}{[(a_{15}^{'})^{(1)}+(a_{15}^{''})^{(1)}(T_{14}^{*})]}$<br>$T_{13}^{*}=\frac{(b_{13})^{(1)}T_{14}^{*}}{[(b_{13}^{'})^{(1)}-(b_{13}^{''})^{(1)}(G^{*})]}$ , $T_{15}^{*}=\frac{(b_{15})^{(1)}T_{14}^{*}}{[(b_{15}^{'})^{(1)}-(b_{15}^{''})^{(1)}(G^{*})]}$<br>Obviously, these values represent an equilibrium solution | |
| Finally we obtain the unique solution | |
| $G_{17}^{*}$ given by $\varphi((G_{19})^{*})=0$ , $T_{17}^{*}$ given by $f(T_{17}^{*})=0$ and | 512 |
| $G_{16}^{*}=\frac{(a_{16})^{(2)}G_{17}^{*}}{[(a_{16}^{'})^{(2)}+(a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18}^{*}=\frac{(a_{18})^{(2)}G_{17}^{*}}{[(a_{18}^{'})^{(2)}+(a_{18}^{''})^{(2)}(T_{17}^{*})]}$ | 513 |
| $T_{16}^{*}=\frac{(b_{16})^{(2)}T_{17}^{*}}{[(b_{16}^{'})^{(2)}-(b_{16}^{''})^{(2)}((G_{19})^{*})]}$ , $T_{18}^{*}=\frac{(b_{18})^{(2)}T_{17}^{*}}{[(b_{18}^{'})^{(2)}-(b_{18}^{''})^{(2)}((G_{19})^{*})]}$ | 514 |
| Obviously, these values represent an equilibrium solution | |
| Finally we obtain the unique solution<br>$G_{21}^{*}$ given by $\varphi((G_{23})^{*})=0$ , $T_{21}^{*}$ given by $f(T_{21}^{*})=0$ and<br>$G_{20}^{*}=\frac{(a_{20})^{(3)}G_{21}^{*}}{[(a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21}^{*})]}$ , $G_{22}^{*}=\frac{(a_{22})^{(3)}G_{21}^{*}}{[(a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21}^{*})]}$<br>$T_{20}^{*}=\frac{(b_{20})^{(3)}T_{21}^{*}}{[(b_{20}^{'})^{(3)}-(b_{20}^{''})^{(3)}(G_{23}{}^{*})]}$ , $T_{22}^{*}=\frac{(b_{22})^{(3)}T_{21}^{*}}{[(b_{22}^{'})^{(3)}-(b_{22}^{''})^{(3)}(G_{23}{}^{*})]}$<br>Obviously, these values represent an equilibrium solution of global equations | 515 |
| Finally we obtain the unique solution<br>$G_{25}^{*}$ given by $\varphi(G_{27})=0$ , $T_{25}^{*}$ given by $f(T_{25}^{*})=0$ and<br>$G_{24}^{*}=\frac{(a_{24})^{(4)}G_{25}^{*}}{[(a_{24}^{'})^{(4)}+(a_{24}^{''})^{(4)}(T_{25}^{*})]}$ , $G_{26}^{*}=\frac{(a_{26})^{(4)}G_{25}^{*}}{[(a_{26}^{'})^{(4)}+(a_{26}^{''})^{(4)}(T_{25}^{*})]}$ | 516 |
| $T_{24}^{*}=\frac{(b_{24})^{(4)}T_{25}^{*}}{[(b_{24}^{'})^{(4)}-(b_{24}^{''})^{(4)}((G_{27})^{*})]}$ , $T_{26}^{*}=\frac{(b_{26})^{(4)}T_{25}^{*}}{[(b_{26}^{'})^{(4)}-(b_{26}^{''})^{(4)}((G_{27})^{*})]}$<br>Obviously, these values represent an equilibrium solution of global equations | 517 |
| Finally we obtain the unique solution<br>$G_{29}^{*}$ given by $\varphi((G_{31})^{*})=0$ , $T_{29}^{*}$ given by $f(T_{29}^{*})=0$ and<br>$G_{28}^{*}=\frac{(a_{28})^{(5)}G_{29}^{*}}{[(a_{28}^{'})^{(5)}+(a_{28}^{''})^{(5)}(T_{29}^{*})]}$ , $G_{30}^{*}=\frac{(a_{30})^{(5)}G_{29}^{*}}{[(a_{30}^{'})^{(5)}+(a_{30}^{''})^{(5)}(T_{29}^{*})]}$ | 518 |

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| $T_{28}^{*} = \frac{(b_{28})^{(5)}T_{29}^{*}}{[(b_{28}^{'})^{(5)}-(b_{28}^{''})^{(5)}((G_{31})^{*})]}$ , $T_{30}^{*} = \frac{(b_{30})^{(5)}T_{29}^{*}}{[(b_{30}^{'})^{(5)}-(b_{30}^{''})^{(5)}((G_{31})^{*})]}$<br>Obviously, these values represent an equilibrium solution of global equations | 519 |
| Finally we obtain the unique solution<br>$G_{33}^{*}$ given by $\varphi((G_{35})^{*}) = 0$ , $T_{33}^{*}$ given by $f(T_{33}^{*}) = 0$ and<br>$G_{32}^{*} = \frac{(a_{32})^{(6)}G_{33}^{*}}{[(a_{32}^{'})^{(6)}+(a_{32}^{''})^{(6)}(T_{33}^{*})]}$ , $G_{34}^{*} = \frac{(a_{34})^{(6)}G_{33}^{*}}{[(a_{34}^{'})^{(6)}+(a_{34}^{''})^{(6)}(T_{33}^{*})]}$ | 520 |
| $T_{32}^{*} = \frac{(b_{32})^{(6)}T_{33}^{*}}{[(b_{32}^{'})^{(6)}-(b_{32}^{''})^{(6)}((G_{35})^{*})]}$ , $T_{34}^{*} = \frac{(b_{34})^{(6)}T_{33}^{*}}{[(b_{34}^{'})^{(6)}-(b_{34}^{''})^{(6)}((G_{35})^{*})]}$<br>Obviously, these values represent an equilibrium solution of global equations | 521 |
| Finally we obtain the unique solution<br>$G_{37}^{*}$ given by $\varphi((G_{39})^{*}) = 0$ , $T_{37}^{*}$ given by $f(T_{37}^{*}) = 0$ and<br>$G_{36}^{*} = \frac{(a_{36})^{(7)}G_{37}^{*}}{[(a_{36}^{'})^{(7)}+(a_{36}^{''})^{(7)}(T_{37}^{*})]}$ , $G_{38}^{*} = \frac{(a_{38})^{(7)}G_{37}^{*}}{[(a_{38}^{'})^{(7)}+(a_{38}^{''})^{(7)}(T_{37}^{*})]}$<br>$T_{36}^{*} = \frac{(b_{36})^{(7)}T_{37}^{*}}{[(b_{36}^{'})^{(7)}-(b_{36}^{''})^{(7)}((G_{39})^{*})]}$ , $T_{38}^{*} = \frac{(b_{38})^{(7)}T_{37}^{*}}{[(b_{38}^{'})^{(7)}-(b_{38}^{''})^{(7)}((G_{39})^{*})]}$ | 522 |
| Finally we obtain the unique solution<br>$G_{41}^{*}$ given by $\varphi((G_{43})^{*}) = 0$ , $T_{41}^{*}$ given by $f(T_{41}^{*}) = 0$ and<br>$G_{40}^{*} = \frac{(a_{40})^{(8)}G_{41}^{*}}{[(a_{40}^{'})^{(8)}+(a_{40}^{''})^{(8)}(T_{41}^{*})]}$ , $G_{42}^{*} = \frac{(a_{42})^{(8)}G_{41}^{*}}{[(a_{42}^{'})^{(8)}+(a_{42}^{''})^{(8)}(T_{41}^{*})]}$<br>$T_{40}^{*} = \frac{(b_{40})^{(8)}T_{41}^{*}}{[(b_{40}^{'})^{(8)}-(b_{40}^{''})^{(8)}((G_{43})^{*})]}$ , $T_{42}^{*} = \frac{(b_{42})^{(8)}T_{41}^{*}}{[(b_{42}^{'})^{(8)}-(b_{42}^{''})^{(8)}((G_{43})^{*})]}$ | 523 |
| Finally we obtain the unique solution of 89 to 99<br>$G_{45}^{*}$ given by $\varphi((G_{47})^{*}) = 0$ , $T_{45}^{*}$ given by $f(T_{45}^{*}) = 0$ and<br>$G_{44}^{*} = \frac{(a_{44})^{(9)}G_{45}^{*}}{[(a_{44}^{'})^{(9)}+(a_{44}^{''})^{(9)}(T_{45}^{*})]}$ , $G_{46}^{*} = \frac{(a_{46})^{(9)}G_{45}^{*}}{[(a_{46}^{'})^{(9)}+(a_{46}^{''})^{(9)}(T_{45}^{*})]}$<br>$T_{44}^{*} = \frac{(b_{44})^{(9)}T_{45}^{*}}{[(b_{44}^{'})^{(9)}-(b_{44}^{''})^{(9)}((G_{47})^{*})]}$ , $T_{46}^{*} = \frac{(b_{46})^{(9)}T_{45}^{*}}{[(b_{46}^{'})^{(9)}-(b_{46}^{''})^{(9)}((G_{47})^{*})]}$ | 523 A |

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| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(1)}$ *and* $(b_i^{''})^{(1)}$ Belong to $C^{(1)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>**Proof:** Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial (a_{14}^{''})^{(1)}}{\partial T_{14}}(T_{14}^*) = (q_{14})^{(1)} \quad , \frac{\partial (b_i^{''})^{(1)}}{\partial G_j}(G^*) = s_{ij}$ | 524 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{13}}{dt} = -\left((a_{13}^{'})^{(1)} + (p_{13})^{(1)}\right)\mathbb{G}_{13} + (a_{13})^{(1)}\mathbb{G}_{14} - (q_{13})^{(1)}G_{13}^*\mathbb{T}_{14}$ | 525 |
| $\frac{d\mathbb{G}_{14}}{dt} = -\left((a_{14}^{'})^{(1)} + (p_{14})^{(1)}\right)\mathbb{G}_{14} + (a_{14})^{(1)}\mathbb{G}_{13} - (q_{14})^{(1)}G_{14}^*\mathbb{T}_{14}$ | 526 |
| $\frac{d\mathbb{G}_{15}}{dt} = -\left((a_{15}^{'})^{(1)} + (p_{15})^{(1)}\right)\mathbb{G}_{15} + (a_{15})^{(1)}\mathbb{G}_{14} - (q_{15})^{(1)}G_{15}^*\mathbb{T}_{14}$ | 527 |
| $\frac{d\mathbb{T}_{13}}{dt} = -\left((b_{13}^{'})^{(1)} - (r_{13})^{(1)}\right)\mathbb{T}_{13} + (b_{13})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(13)(j)}T_{13}^*\mathbb{G}_j\right)$ | 528 |
| $\frac{d\mathbb{T}_{14}}{dt} = -\left((b_{14}^{'})^{(1)} - (r_{14})^{(1)}\right)\mathbb{T}_{14} + (b_{14})^{(1)}\mathbb{T}_{13} + \sum_{j=13}^{15}\left(s_{(14)(j)}T_{14}^*\mathbb{G}_j\right)$ | 529 |
| $\frac{d\mathbb{T}_{15}}{dt} = -\left((b_{15}^{'})^{(1)} - (r_{15})^{(1)}\right)\mathbb{T}_{15} + (b_{15})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(15)(j)}T_{15}^*\mathbb{G}_j\right)$ | 530 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}$ Belong to $C^{(2)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable | 531 |
| Proof: Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- | |
| $G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$ | 532 |
| $\frac{\partial (a_{17}^{''})^{(2)}}{\partial T_{17}}(T_{17}^*) = (q_{17})^{(2)} \quad , \frac{\partial (b_i^{''})^{(2)}}{\partial G_j}((G_{19})^*) = s_{ij}$ | 533 |
| taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{16}}{dt} = -\left((a_{16}^{'})^{(2)} + (p_{16})^{(2)}\right)\mathbb{G}_{16} + (a_{16})^{(2)}\mathbb{G}_{17} - (q_{16})^{(2)}G_{16}^*\mathbb{T}_{17}$ | 534 |
| $\frac{d\mathbb{G}_{17}}{dt} = -\left((a_{17}^{'})^{(2)} + (p_{17})^{(2)}\right)\mathbb{G}_{17} + (a_{17})^{(2)}\mathbb{G}_{16} - (q_{17})^{(2)}G_{17}^*\mathbb{T}_{17}$ | 535 |

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| $\frac{d\mathbb{G}_{18}}{dt} = -\left((a_{18}^{'})^{(2)} + (p_{18})^{(2)}\right)\mathbb{G}_{18} + (a_{18})^{(2)}\mathbb{G}_{17} - (q_{18})^{(2)}G_{18}^{*}\mathbb{T}_{17}$ | 536 |
| $\frac{d\mathbb{T}_{16}}{dt} = -\left((b_{16}^{'})^{(2)} - (r_{16})^{(2)}\right)\mathbb{T}_{16} + (b_{16})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(16)(j)}T_{16}^{*}\mathbb{G}_{j}\right)$ | 537 |
| $\frac{d\mathbb{T}_{17}}{dt} = -\left((b_{17}^{'})^{(2)} - (r_{17})^{(2)}\right)\mathbb{T}_{17} + (b_{17})^{(2)}\mathbb{T}_{16} + \sum_{j=16}^{18}\left(s_{(17)(j)}T_{17}^{*}\mathbb{G}_{j}\right)$ | 538 |
| $\frac{d\mathbb{T}_{18}}{dt} = -\left((b_{18}^{'})^{(2)} - (r_{18})^{(2)}\right)\mathbb{T}_{18} + (b_{18})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(18)(j)}T_{18}^{*}\mathbb{G}_{j}\right)$ | 539 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(3)}$ $and$ $(b_i^{''})^{(3)}$ Belong to $C^{(3)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial(a_{21}^{''})^{(3)}}{\partial T_{21}}(T_{21}^{*}) = (q_{21})^{(3)}$ , $\frac{\partial(b_i^{''})^{(3)}}{\partial G_j}((G_{23})^{*}) = s_{ij}$ | 540 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{20}}{dt} = -\left((a_{20}^{'})^{(3)} + (p_{20})^{(3)}\right)\mathbb{G}_{20} + (a_{20})^{(3)}\mathbb{G}_{21} - (q_{20})^{(3)}G_{20}^{*}\mathbb{T}_{21}$ | 541 |
| $\frac{d\mathbb{G}_{21}}{dt} = -\left((a_{21}^{'})^{(3)} + (p_{21})^{(3)}\right)\mathbb{G}_{21} + (a_{21})^{(3)}\mathbb{G}_{20} - (q_{21})^{(3)}G_{21}^{*}\mathbb{T}_{21}$ | 542 |
| $\frac{d\mathbb{G}_{22}}{dt} = -\left((a_{22}^{'})^{(3)} + (p_{22})^{(3)}\right)\mathbb{G}_{22} + (a_{22})^{(3)}\mathbb{G}_{21} - (q_{22})^{(3)}G_{22}^{*}\mathbb{T}_{21}$ | 543 |
| $\frac{d\mathbb{T}_{20}}{dt} = -\left((b_{20}^{'})^{(3)} - (r_{20})^{(3)}\right)\mathbb{T}_{20} + (b_{20})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(20)(j)}T_{20}^{*}\mathbb{G}_{j}\right)$ | 544 |
| $\frac{d\mathbb{T}_{21}}{dt} = -\left((b_{21}^{'})^{(3)} - (r_{21})^{(3)}\right)\mathbb{T}_{21} + (b_{21})^{(3)}\mathbb{T}_{20} + \sum_{j=20}^{22}\left(s_{(21)(j)}T_{21}^{*}\mathbb{G}_{j}\right)$ | 545 |
| $\frac{d\mathbb{T}_{22}}{dt} = -\left((b_{22}^{'})^{(3)} - (r_{22})^{(3)}\right)\mathbb{T}_{22} + (b_{22})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(22)(j)}T_{22}^{*}\mathbb{G}_{j}\right)$ | 546 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(4)}$ $and$ $(b_i^{''})^{(4)}$ Belong to $C^{(4)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 547 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$ | 548 |

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| $\frac{\partial (a_{25}^{''})^{(4)}}{\partial T_{25}}(T_{25}^{*}) = (q_{25})^{(4)}$ , $\frac{\partial (b_i^{''})^{(4)}}{\partial G_j}((G_{27})^{*}) = s_{ij}$ | |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{24}}{dt} = -\left((a_{24}^{'})^{(4)} + (p_{24})^{(4)}\right)\mathbb{G}_{24} + (a_{24})^{(4)}\mathbb{G}_{25} - (q_{24})^{(4)}G_{24}^{*}\mathbb{T}_{25}$ | 549 |
| $\frac{d\mathbb{G}_{25}}{dt} = -\left((a_{25}^{'})^{(4)} + (p_{25})^{(4)}\right)\mathbb{G}_{25} + (a_{25})^{(4)}\mathbb{G}_{24} - (q_{25})^{(4)}G_{25}^{*}\mathbb{T}_{25}$ | 550 |
| $\frac{d\mathbb{G}_{26}}{dt} = -\left((a_{26}^{'})^{(4)} + (p_{26})^{(4)}\right)\mathbb{G}_{26} + (a_{26})^{(4)}\mathbb{G}_{25} - (q_{26})^{(4)}G_{26}^{*}\mathbb{T}_{25}$ | 551 |
| $\frac{d\mathbb{T}_{24}}{dt} = -\left((b_{24}^{'})^{(4)} - (r_{24})^{(4)}\right)\mathbb{T}_{24} + (b_{24})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(24)(j)}T_{24}^{*}\mathbb{G}_j\right)$ | 552 |
| $\frac{d\mathbb{T}_{25}}{dt} = -\left((b_{25}^{'})^{(4)} - (r_{25})^{(4)}\right)\mathbb{T}_{25} + (b_{25})^{(4)}\mathbb{T}_{24} + \sum_{j=24}^{26}\left(s_{(25)(j)}T_{25}^{*}\mathbb{G}_j\right)$ | 553 |
| $\frac{d\mathbb{T}_{26}}{dt} = -\left((b_{26}^{'})^{(4)} - (r_{26})^{(4)}\right)\mathbb{T}_{26} + (b_{26})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(26)(j)}T_{26}^{*}\mathbb{G}_j\right)$ | 554 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 5:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(5)}$ *and* $(b_i^{''})^{(5)}$ Belong to $C^{(5)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 555 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^{*} + \mathbb{G}_i \quad , T_i = T_i^{*} + \mathbb{T}_i$<br><br>$\frac{\partial (a_{29}^{''})^{(5)}}{\partial T_{29}}(T_{29}^{*}) = (q_{29})^{(5)}$ , $\frac{\partial (b_i^{''})^{(5)}}{\partial G_j}((G_{31})^{*}) = s_{ij}$ | 556 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{28}}{dt} = -\left((a_{28}^{'})^{(5)} + (p_{28})^{(5)}\right)\mathbb{G}_{28} + (a_{28})^{(5)}\mathbb{G}_{29} - (q_{28})^{(5)}G_{28}^{*}\mathbb{T}_{29}$ | 557 |
| $\frac{d\mathbb{G}_{29}}{dt} = -\left((a_{29}^{'})^{(5)} + (p_{29})^{(5)}\right)\mathbb{G}_{29} + (a_{29})^{(5)}\mathbb{G}_{28} - (q_{29})^{(5)}G_{29}^{*}\mathbb{T}_{29}$ | 558 |
| $\frac{d\mathbb{G}_{30}}{dt} = -\left((a_{30}^{'})^{(5)} + (p_{30})^{(5)}\right)\mathbb{G}_{30} + (a_{30})^{(5)}\mathbb{G}_{29} - (q_{30})^{(5)}G_{30}^{*}\mathbb{T}_{29}$ | 559 |
| $\frac{d\mathbb{T}_{28}}{dt} = -\left((b_{28}^{'})^{(5)} - (r_{28})^{(5)}\right)\mathbb{T}_{28} + (b_{28})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(28)(j)}T_{28}^{*}\mathbb{G}_j\right)$ | 560 |
| $\frac{d\mathbb{T}_{29}}{dt} = -\left((b_{29}^{'})^{(5)} - (r_{29})^{(5)}\right)\mathbb{T}_{29} + (b_{29})^{(5)}\mathbb{T}_{28} + \sum_{j=28}^{30}\left(s_{(29)(j)}T_{29}^{*}\mathbb{G}_j\right)$ | 561 |
| $\frac{d\mathbb{T}_{30}}{dt} = -\left((b_{30}^{'})^{(5)} - (r_{30})^{(5)}\right)\mathbb{T}_{30} + (b_{30})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(30)(j)}T_{30}^{*}\mathbb{G}_j\right)$ | 562 |
| **ASYMPTOTIC STABILITY ANALYSIS** | 563 |

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| **Theorem 6:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(6)}$ $and$ $(b_i^{''})^{(6)}$ Belong to $C^{(6)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial (a_{33}^{''})^{(6)}}{\partial T_{33}}(T_{33}^*) = (q_{33})^{(6)}$ , $\frac{\partial (b_i^{''})^{(6)}}{\partial G_j}((G_{35})^*) = s_{ij}$ | 564 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{32}}{dt} = -\left((a_{32}^{'})^{(6)} + (p_{32})^{(6)}\right)\mathbb{G}_{32} + (a_{32})^{(6)}\mathbb{G}_{33} - (q_{32})^{(6)}G_{32}^*\mathbb{T}_{33}$ | 565 |
| $\frac{d\mathbb{G}_{33}}{dt} = -\left((a_{33}^{'})^{(6)} + (p_{33})^{(6)}\right)\mathbb{G}_{33} + (a_{33})^{(6)}\mathbb{G}_{32} - (q_{33})^{(6)}G_{33}^*\mathbb{T}_{33}$ | 566 |
| $\frac{d\mathbb{G}_{34}}{dt} = -\left((a_{34}^{'})^{(6)} + (p_{34})^{(6)}\right)\mathbb{G}_{34} + (a_{34})^{(6)}\mathbb{G}_{33} - (q_{34})^{(6)}G_{34}^*\mathbb{T}_{33}$ | 567 |
| $\frac{d\mathbb{T}_{32}}{dt} = -\left((b_{32}^{'})^{(6)} - (r_{32})^{(6)}\right)\mathbb{T}_{32} + (b_{32})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(32)(j)}T_{32}^*\mathbb{G}_j\right)$ | 568 |
| $\frac{d\mathbb{T}_{33}}{dt} = -\left((b_{33}^{'})^{(6)} - (r_{33})^{(6)}\right)\mathbb{T}_{33} + (b_{33})^{(6)}\mathbb{T}_{32} + \sum_{j=32}^{34}\left(s_{(33)(j)}T_{33}^*\mathbb{G}_j\right)$ | 569 |
| $\frac{d\mathbb{T}_{34}}{dt} = -\left((b_{34}^{'})^{(6)} - (r_{34})^{(6)}\right)\mathbb{T}_{34} + (b_{34})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(34)(j)}T_{34}^*\mathbb{G}_j\right)$ | 570 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 7:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(7)}$ $and$ $(b_i^{''})^{(7)}$ Belong to $C^{(7)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 571 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial (a_{37}^{''})^{(7)}}{\partial T_{37}}(T_{37}^*) = (q_{37})^{(7)}$ , $\frac{\partial (b_i^{''})^{(7)}}{\partial G_j}((G_{39})^{**}) = s_{ij}$ | 572 |
| Then taking into account equations and neglecting the terms of power 2, we obtain from | |
| $\frac{d\mathbb{G}_{36}}{dt} = -\left((a_{36}^{'})^{(7)} + (p_{36})^{(7)}\right)\mathbb{G}_{36} + (a_{36})^{(7)}\mathbb{G}_{37} - (q_{36})^{(7)}G_{36}^*\mathbb{T}_{37}$ | 573 |
| $\frac{d\mathbb{G}_{37}}{dt} = -\left((a_{37}^{'})^{(7)} + (p_{37})^{(7)}\right)\mathbb{G}_{37} + (a_{37})^{(7)}\mathbb{G}_{36} - (q_{37})^{(7)}G_{37}^*\mathbb{T}_{37}$ | 574 |

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| $\frac{d\mathbb{G}_{38}}{dt} = -\left((a'_{38})^{(7)} + (p_{38})^{(7)}\right)\mathbb{G}_{38} + (a_{38})^{(7)}\mathbb{G}_{37} - (q_{38})^{(7)}G^*_{38}\mathbb{T}_{37}$ | 575 |
| $\frac{d\mathbb{T}_{36}}{dt} = -\left((b'_{36})^{(7)} - (r_{36})^{(7)}\right)\mathbb{T}_{36} + (b_{36})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(36)(j)}T^*_{36}\mathbb{G}_j\right)$ | 576 |
| $\frac{d\mathbb{T}_{37}}{dt} = -\left((b'_{37})^{(7)} - (r_{37})^{(7)}\right)\mathbb{T}_{37} + (b_{37})^{(7)}\mathbb{T}_{36} + \sum_{j=36}^{38}\left(s_{(37)(j)}T^*_{37}\mathbb{G}_j\right)$ | 578 |
| $\frac{d\mathbb{T}_{38}}{dt} = -\left((b'_{38})^{(7)} - (r_{38})^{(7)}\right)\mathbb{T}_{38} + (b_{38})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(38)(j)}T^*_{38}\mathbb{G}_j\right)$ | 579 |
| Obviously, these values represent an equilibrium solution<br><br>**ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 8:** If the conditions of the previous theorem are satisfied and if the functions $(a_i'')^{(8)}$ *and* $(b_i'')^{(8)}$ Belong to $C^{(8)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial(a''_{41})^{(8)}}{\partial T_{41}}(T^*_{41}) = (q_{41})^{(8)}$ , $\frac{\partial(b_i'')^{(8)}}{\partial G_j}((G_{43})^*) = s_{ij}$ | 580 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{40}}{dt} = -\left((a'_{40})^{(8)} + (p_{40})^{(8)}\right)\mathbb{G}_{40} + (a_{40})^{(8)}\mathbb{G}_{41} - (q_{40})^{(8)}G^*_{40}\mathbb{T}_{41}$ | 581 |
| $\frac{d\mathbb{G}_{41}}{dt} = -\left((a'_{41})^{(8)} + (p_{41})^{(8)}\right)\mathbb{G}_{41} + (a_{41})^{(8)}\mathbb{G}_{40} - (q_{41})^{(8)}G^*_{41}\mathbb{T}_{41}$ | 582 |
| $\frac{d\mathbb{G}_{42}}{dt} = -\left((a'_{42})^{(8)} + (p_{42})^{(8)}\right)\mathbb{G}_{42} + (a_{42})^{(8)}\mathbb{G}_{41} - (q_{42})^{(8)}G^*_{42}\mathbb{T}_{41}$ | 583 |
| $\frac{d\mathbb{T}_{40}}{dt} = -\left((b'_{40})^{(8)} - (r_{40})^{(8)}\right)\mathbb{T}_{40} + (b_{40})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(40)(j)}T^*_{40}\mathbb{G}_j\right)$ | 584 |
| $\frac{d\mathbb{T}_{41}}{dt} = -\left((b'_{41})^{(8)} - (r_{41})^{(8)}\right)\mathbb{T}_{41} + (b_{41})^{(8)}\mathbb{T}_{40} + \sum_{j=40}^{42}\left(s_{(41)(j)}T^*_{41}\mathbb{G}_j\right)$ | 585 |
| $\frac{d\mathbb{T}_{42}}{dt} = -\left((b'_{42})^{(8)} - (r_{42})^{(8)}\right)\mathbb{T}_{42} + (b_{42})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(42)(j)}T^*_{42}\mathbb{G}_j\right)$ | 586 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 9:** If the conditions of the previous theorem are satisfied and if the functions $(a_i'')^{(9)}$ *and* $(b_i'')^{(9)}$ Belong to $C^{(9)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br>**Proof: Denote** | 586 A |

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|---|---|
| | |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br>$\frac{\partial (a_{45}^{''})^{(9)}}{\partial T_{45}}(T_{45}^*) = (q_{45})^{(9)} \ , \ \frac{\partial (b_i^{''})^{(9)}}{\partial G_j}((G_{47})^*) = s_{ij}$ | |
| Then taking into account equations 89 to 99 and neglecting the terms of power 2, we obtain from 99 to 44 | |
| $\frac{d\mathbb{G}_{44}}{dt} = -\left((a_{44}^{'})^{(9)} + (p_{44})^{(9)}\right)\mathbb{G}_{44} + (a_{44})^{(9)}\mathbb{G}_{45} - (q_{44})^{(9)}G_{44}^*\mathbb{T}_{45}$ | 586 B |
| $\frac{d\mathbb{G}_{45}}{dt} = -\left((a_{45}^{'})^{(9)} + (p_{45})^{(9)}\right)\mathbb{G}_{45} + (a_{45})^{(9)}\mathbb{G}_{44} - (q_{45})^{(9)}G_{45}^*\mathbb{T}_{45}$ | 586 C |
| $\frac{d\mathbb{G}_{46}}{dt} = -\left((a_{46}^{'})^{(9)} + (p_{46})^{(9)}\right)\mathbb{G}_{46} + (a_{46})^{(9)}\mathbb{G}_{45} - (q_{46})^{(9)}G_{46}^*\mathbb{T}_{45}$ | 586 D |
| $\frac{d\mathbb{T}_{44}}{dt} = -\left((b_{44}^{'})^{(9)} - (r_{44})^{(9)}\right)\mathbb{T}_{44} + (b_{44})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(44)(j)}T_{44}^*\mathbb{G}_j\right)$ | 586 E |
| $\frac{d\mathbb{T}_{45}}{dt} = -\left((b_{45}^{'})^{(9)} - (r_{45})^{(9)}\right)\mathbb{T}_{45} + (b_{45})^{(9)}\mathbb{T}_{44} + \sum_{j=44}^{46}\left(s_{(45)(j)}T_{45}^*\mathbb{G}_j\right)$ | 586 F |
| $\frac{d\mathbb{T}_{46}}{dt} = -\left((b_{46}^{'})^{(9)} - (r_{46})^{(9)}\right)\mathbb{T}_{46} + (b_{46})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(46)(j)}T_{46}^*\mathbb{G}_j\right)$ | 586 G |
| **The characteristic equation of this system is** | 587 |
| $\left((\lambda)^{(1)} + (b_{15}^{'})^{(1)} - (r_{15})^{(1)}\right)\{\left((\lambda)^{(1)} + (a_{15}^{'})^{(1)} + (p_{15})^{(1)}\right)$<br>$\left[\left((\lambda)^{(1)} + (a_{13}^{'})^{(1)} + (p_{13})^{(1)}\right)(q_{14})^{(1)}G_{14}^* + (a_{14})^{(1)}(q_{13})^{(1)}G_{13}^*\right)\Big]$<br>$\left((\lambda)^{(1)} + (b_{13}^{'})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(14)}T_{14}^* + (b_{14})^{(1)}s_{(13),(14)}T_{14}^*\Big)$<br>$+\left(\left((\lambda)^{(1)} + (a_{14}^{'})^{(1)} + (p_{14})^{(1)}\right)(q_{13})^{(1)}G_{13}^* + (a_{13})^{(1)}(q_{14})^{(1)}G_{14}^*\right)$<br>$\left((\lambda)^{(1)} + (b_{13}^{'})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(13)}T_{14}^* + (b_{14})^{(1)}s_{(13),(13)}T_{13}^*\Big)$<br>$\left((\lambda)^{(1)}\right)^2 + \left((a_{13}^{'})^{(1)} + (a_{14}^{'})^{(1)} + (p_{13})^{(1)} + (p_{14})^{(1)}\right)(\lambda)^{(1)}\Big)$<br>$\left((\lambda)^{(1)}\right)^2 + \left((b_{13}^{'})^{(1)} + (b_{14}^{'})^{(1)} - (r_{13})^{(1)} + (r_{14})^{(1)}\right)(\lambda)^{(1)}\Big)$<br>$+\left(\left((\lambda)^{(1)}\right)^2 + \left((a_{13}^{'})^{(1)} + (a_{14}^{'})^{(1)} + (p_{13})^{(1)} + (p_{14})^{(1)}\right)(\lambda)^{(1)}\right)(q_{15})^{(1)}G_{15}$<br>$+\left((\lambda)^{(1)} + (a_{13}^{'})^{(1)} + (p_{13})^{(1)}\right)\left((a_{15})^{(1)}(q_{14})^{(1)}G_{14}^* + (a_{14})^{(1)}(a_{15})^{(1)}(q_{13})^{(1)}G_{13}^*\right)$<br>$\left((\lambda)^{(1)} + (b_{13}^{'})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(15)}T_{14}^* + (b_{14})^{(1)}s_{(13),(15)}T_{13}^*\Big)\} = 0$ | |

$$+$$

$$((\lambda)^{(2)} + (b'_{18})^{(2)} - (r_{18})^{(2)})\{((\lambda)^{(2)} + (a'_{18})^{(2)} + (p_{18})^{(2)})$$

$$\left[\left(((\lambda)^{(2)} + (a'_{16})^{(2)} + (p_{16})^{(2)})(q_{17})^{(2)}\mathrm{G}^{*}_{17} + (a_{17})^{(2)}(q_{16})^{(2)}\mathrm{G}^{*}_{16}\right)\right]$$

$$\left(((\lambda)^{(2)} + (b'_{16})^{(2)} - (r_{16})^{(2)})s_{(17),(17)}\mathrm{T}^{*}_{17} + (b_{17})^{(2)}s_{(16),(17)}\mathrm{T}^{*}_{17}\right)$$

$$+\left(((\lambda)^{(2)} + (a'_{17})^{(2)} + (p_{17})^{(2)})(q_{16})^{(2)}\mathrm{G}^{*}_{16} + (a_{16})^{(2)}(q_{17})^{(2)}\mathrm{G}^{*}_{17}\right)$$

$$\left(((\lambda)^{(2)} + (b'_{16})^{(2)} - (r_{16})^{(2)})s_{(17),(16)}\mathrm{T}^{*}_{17} + (b_{17})^{(2)}s_{(16),(16)}\mathrm{T}^{*}_{16}\right)$$

$$\left((\lambda)^{(2)}\right)^{2} + \left((a'_{16})^{(2)} + (a'_{17})^{(2)} + (p_{16})^{(2)} + (p_{17})^{(2)}\right)(\lambda)^{(2)}\Big)$$

$$\left((\lambda)^{(2)}\right)^{2} + \left((b'_{16})^{(2)} + (b'_{17})^{(2)} - (r_{16})^{(2)} + (r_{17})^{(2)}\right)(\lambda)^{(2)}\Big)$$

$$+\left(\left((\lambda)^{(2)}\right)^{2} + \left((a'_{16})^{(2)} + (a'_{17})^{(2)} + (p_{16})^{(2)} + (p_{17})^{(2)}\right)(\lambda)^{(2)}\right)(q_{18})^{(2)}\mathrm{G}_{18}$$

$$+\left((\lambda)^{(2)} + (a'_{16})^{(2)} + (p_{16})^{(2)}\right)\left((a_{18})^{(2)}(q_{17})^{(2)}\mathrm{G}^{*}_{17} + (a_{17})^{(2)}(a_{18})^{(2)}(q_{16})^{(2)}\mathrm{G}^{*}_{16}\right)$$

$$\left(((\lambda)^{(2)} + (b'_{16})^{(2)} - (r_{16})^{(2)})s_{(17),(18)}\mathrm{T}^{*}_{17} + (b_{17})^{(2)}s_{(16),(18)}\mathrm{T}^{*}_{16}\right)\} = 0$$

$$+$$

$$((\lambda)^{(3)} + (b'_{22})^{(3)} - (r_{22})^{(3)})\{((\lambda)^{(3)} + (a'_{22})^{(3)} + (p_{22})^{(3)})$$

$$\left[\left(((\lambda)^{(3)} + (a'_{20})^{(3)} + (p_{20})^{(3)})(q_{21})^{(3)}G^{*}_{21} + (a_{21})^{(3)}(q_{20})^{(3)}G^{*}_{20}\right)\right]$$

$$\left(((\lambda)^{(3)} + (b'_{20})^{(3)} - (r_{20})^{(3)})s_{(21),(21)}T^{*}_{21} + (b_{21})^{(3)}s_{(20),(21)}T^{*}_{21}\right)$$

$$+\left(((\lambda)^{(3)} + (a'_{21})^{(3)} + (p_{21})^{(3)})(q_{20})^{(3)}G^{*}_{20} + (a_{20})^{(3)}(q_{21})^{(1)}G^{*}_{21}\right)$$

$$\left(((\lambda)^{(3)} + (b'_{20})^{(3)} - (r_{20})^{(3)})s_{(21),(20)}T^{*}_{21} + (b_{21})^{(3)}s_{(20),(20)}T^{*}_{20}\right)$$

$$\left((\lambda)^{(3)}\right)^{2} + \left((a'_{20})^{(3)} + (a'_{21})^{(3)} + (p_{20})^{(3)} + (p_{21})^{(3)}\right)(\lambda)^{(3)}\Big)$$

$$\left((\lambda)^{(3)}\right)^{2} + \left((b'_{20})^{(3)} + (b'_{21})^{(3)} - (r_{20})^{(3)} + (r_{21})^{(3)}\right)(\lambda)^{(3)}\Big)$$

$$+\left(\left((\lambda)^{(3)}\right)^{2} + \left((a'_{20})^{(3)} + (a'_{21})^{(3)} + (p_{20})^{(3)} + (p_{21})^{(3)}\right)(\lambda)^{(3)}\right)(q_{22})^{(3)}G_{22}$$

$$+\left((\lambda)^{(3)} + (a'_{20})^{(3)} + (p_{20})^{(3)}\right)\left((a_{22})^{(3)}(q_{21})^{(3)}G^{*}_{21} + (a_{21})^{(3)}(a_{22})^{(3)}(q_{20})^{(3)}G^{*}_{20}\right)$$

$$\left(((\lambda)^{(3)} + (b'_{20})^{(3)} - (r_{20})^{(3)})s_{(21),(22)}T^{*}_{21} + (b_{21})^{(3)}s_{(20),(22)}T^{*}_{20}\right)\} = 0$$

$$+$$

$$\left((\lambda)^{(4)}+(b'_{26})^{(4)}-(r_{26})^{(4)}\right)\{\left((\lambda)^{(4)}+(a'_{26})^{(4)}+(p_{26})^{(4)}\right)$$

$$\left[\left(\left((\lambda)^{(4)}+(a'_{24})^{(4)}+(p_{24})^{(4)}\right)(q_{25})^{(4)}G^{*}_{25}+(a_{25})^{(4)}(q_{24})^{(4)}G^{*}_{24}\right)\right]$$

$$\left(\left((\lambda)^{(4)}+(b'_{24})^{(4)}-(r_{24})^{(4)}\right)s_{(25),(25)}T^{*}_{25}+(b_{25})^{(4)}s_{(24),(25)}T^{*}_{25}\right)$$

$$+\left(\left((\lambda)^{(4)}+(a'_{25})^{(4)}+(p_{25})^{(4)}\right)(q_{24})^{(4)}G^{*}_{24}+(a_{24})^{(4)}(q_{25})^{(4)}G^{*}_{25}\right)$$

$$\left(\left((\lambda)^{(4)}+(b'_{24})^{(4)}-(r_{24})^{(4)}\right)s_{(25),(24)}T^{*}_{25}+(b_{25})^{(4)}s_{(24),(24)}T^{*}_{24}\right)$$

$$\left(\left((\lambda)^{(4)}\right)^{2}+\left((a'_{24})^{(4)}+(a'_{25})^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)$$

$$\left(\left((\lambda)^{(4)}\right)^{2}+\left((b'_{24})^{(4)}+(b'_{25})^{(4)}-(r_{24})^{(4)}+(r_{25})^{(4)}\right)(\lambda)^{(4)}\right)$$

$$+\left(\left((\lambda)^{(4)}\right)^{2}+\left((a'_{24})^{(4)}+(a'_{25})^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)(q_{26})^{(4)}G_{26}$$

$$+\left((\lambda)^{(4)}+(a'_{24})^{(4)}+(p_{24})^{(4)}\right)\left((a_{26})^{(4)}(q_{25})^{(4)}G^{*}_{25}+(a_{25})^{(4)}(a_{26})^{(4)}(q_{24})^{(4)}G^{*}_{24}\right)$$

$$\left(\left((\lambda)^{(4)}+(b'_{24})^{(4)}-(r_{24})^{(4)}\right)s_{(25),(26)}T^{*}_{25}+(b_{25})^{(4)}s_{(24),(26)}T^{*}_{24}\right)\}=0$$

$$+$$

$$\left((\lambda)^{(5)}+(b'_{30})^{(5)}-(r_{30})^{(5)}\right)\{\left((\lambda)^{(5)}+(a'_{30})^{(5)}+(p_{30})^{(5)}\right)$$

$$\left[\left(\left((\lambda)^{(5)}+(a'_{28})^{(5)}+(p_{28})^{(5)}\right)(q_{29})^{(5)}G^{*}_{29}+(a_{29})^{(5)}(q_{28})^{(5)}G^{*}_{28}\right)\right]$$

$$\left(\left((\lambda)^{(5)}+(b'_{28})^{(5)}-(r_{28})^{(5)}\right)s_{(29),(29)}T^{*}_{29}+(b_{29})^{(5)}s_{(28),(29)}T^{*}_{29}\right)$$

$$+\left(\left((\lambda)^{(5)}+(a'_{29})^{(5)}+(p_{29})^{(5)}\right)(q_{28})^{(5)}G^{*}_{28}+(a_{28})^{(5)}(q_{29})^{(5)}G^{*}_{29}\right)$$

$$\left(\left((\lambda)^{(5)}+(b'_{28})^{(5)}-(r_{28})^{(5)}\right)s_{(29),(28)}T^{*}_{29}+(b_{29})^{(5)}s_{(28),(28)}T^{*}_{28}\right)$$

$$\left(\left((\lambda)^{(5)}\right)^{2}+\left((a'_{28})^{(5)}+(a'_{29})^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)$$

$$\left(\left((\lambda)^{(5)}\right)^{2}+\left((b'_{28})^{(5)}+(b'_{29})^{(5)}-(r_{28})^{(5)}+(r_{29})^{(5)}\right)(\lambda)^{(5)}\right)$$

$$+\left(\left((\lambda)^{(5)}\right)^{2}+\left((a'_{28})^{(5)}+(a'_{29})^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)(q_{30})^{(5)}G_{30}$$

$$+\left((\lambda)^{(5)}+(a'_{28})^{(5)}+(p_{28})^{(5)}\right)\left((a_{30})^{(5)}(q_{29})^{(5)}G^{*}_{29}+(a_{29})^{(5)}(a_{30})^{(5)}(q_{28})^{(5)}G^{*}_{28}\right)$$

$$\left(\left((\lambda)^{(5)}+(b'_{28})^{(5)}-(r_{28})^{(5)}\right)s_{(29),(30)}T^{*}_{29}+(b_{29})^{(5)}s_{(28),(30)}T^{*}_{28}\right)\}=0$$

$$+$$

$$\left((\lambda)^{(6)}+(b_{34}^{'})^{(6)}-(r_{34})^{(6)}\right)\{\left((\lambda)^{(6)}+(a_{34}^{'})^{(6)}+(p_{34})^{(6)}\right)$$

$$\left[\left(\left((\lambda)^{(6)}+(a_{32}^{'})^{(6)}+(p_{32})^{(6)}\right)(q_{33})^{(6)}G_{33}^{*}+(a_{33})^{(6)}(q_{32})^{(6)}G_{32}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(6)}+(b_{32}^{'})^{(6)}-(r_{32})^{(6)}\right)s_{(33),(33)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(33)}T_{33}^{*}\right)$$

$$+\left(\left((\lambda)^{(6)}+(a_{33}^{'})^{(6)}+(p_{33})^{(6)}\right)(q_{32})^{(6)}G_{32}^{*}+(a_{32})^{(6)}(q_{33})^{(6)}G_{33}^{*}\right)$$

$$\left(\left((\lambda)^{(6)}+(b_{32}^{'})^{(6)}-(r_{32})^{(6)}\right)s_{(33),(32)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(32)}T_{32}^{*}\right)$$

$$\left(\left((\lambda)^{(6)}\right)^{2}+\left((a_{32}^{'})^{(6)}+(a_{33}^{'})^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\right)$$

$$\left(\left((\lambda)^{(6)}\right)^{2}+\left((b_{32}^{'})^{(6)}+(b_{33}^{'})^{(6)}-(r_{32})^{(6)}+(r_{33})^{(6)}\right)(\lambda)^{(6)}\right)$$

$$+\left(\left((\lambda)^{(6)}\right)^{2}+\left((a_{32}^{'})^{(6)}+(a_{33}^{'})^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\right)(q_{34})^{(6)}G_{34}$$

$$+\left((\lambda)^{(6)}+(a_{32}^{'})^{(6)}+(p_{32})^{(6)}\right)\left((a_{34})^{(6)}(q_{33})^{(6)}G_{33}^{*}+(a_{33})^{(6)}(a_{34})^{(6)}(q_{32})^{(6)}G_{32}^{*}\right)$$

$$\left(\left((\lambda)^{(6)}+(b_{32}^{'})^{(6)}-(r_{32})^{(6)}\right)s_{(33),(34)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(34)}T_{32}^{*}\right)\}=0$$

$$+$$

$$\left((\lambda)^{(7)}+(b_{38}^{'})^{(7)}-(r_{38})^{(7)}\right)\{\left((\lambda)^{(7)}+(a_{38}^{'})^{(7)}+(p_{38})^{(7)}\right)$$

$$\left[\left(\left((\lambda)^{(7)}+(a_{36}^{'})^{(7)}+(p_{36})^{(7)}\right)(q_{37})^{(7)}G_{37}^{*}+(a_{37})^{(7)}(q_{36})^{(7)}G_{36}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(7)}+(b_{36}^{'})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(37)}T_{37}^{*}+(b_{37})^{(7)}s_{(36),(37)}T_{37}^{*}\right)$$

$$+\left(\left((\lambda)^{(7)}+(a_{37}^{'})^{(7)}+(p_{37})^{(7)}\right)(q_{36})^{(7)}G_{36}^{*}+(a_{36})^{(7)}(q_{37})^{(7)}G_{37}^{*}\right)$$

$$\left(\left((\lambda)^{(7)}+(b_{36}^{'})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(36)}T_{37}^{*}+(b_{37})^{(7)}s_{(36),(36)}T_{36}^{*}\right)$$

$$\left(\left((\lambda)^{(7)}\right)^{2}+\left((a_{36}^{'})^{(7)}+(a_{37}^{'})^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\right)$$

$$\left(\left((\lambda)^{(7)}\right)^{2}+\left((b_{36}^{'})^{(7)}+(b_{37}^{'})^{(7)}-(r_{36})^{(7)}+(r_{37})^{(7)}\right)(\lambda)^{(7)}\right)$$

$$+\left(\left((\lambda)^{(7)}\right)^{2}+\left((a_{36}^{'})^{(7)}+(a_{37}^{'})^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\right)(q_{38})^{(7)}G_{38}$$

$$+\left((\lambda)^{(7)}+(a_{36}^{'})^{(7)}+(p_{36})^{(7)}\right)\left((a_{38})^{(7)}(q_{37})^{(7)}G_{37}^{*}+(a_{37})^{(7)}(a_{38})^{(7)}(q_{36})^{(7)}G_{36}^{*}\right)$$

$$\left(\left((\lambda)^{(7)}+(b_{36}^{'})^{(7)}-(r_{36})^{(7)}\right)s_{(37),(38)}T_{37}^{*}+(b_{37})^{(7)}s_{(36),(38)}T_{36}^{*}\right)\}=0$$

$$+$$

$$\left((\lambda)^{(8)}+(b'_{42})^{(8)}-(r_{42})^{(8)}\right)\{\left((\lambda)^{(8)}+(a'_{42})^{(8)}+(p_{42})^{(8)}\right)$$

$$\left[\left(\left((\lambda)^{(8)}+(a'_{40})^{(8)}+(p_{40})^{(8)}\right)(q_{41})^{(8)}G^{*}_{41}+(a_{41})^{(8)}(q_{40})^{(8)}G^{*}_{40}\right)\right]$$

$$\left(\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(41)}T^{*}_{41}+(b_{41})^{(8)}s_{(40),(41)}T^{*}_{41}\right)$$

$$+\left(\left((\lambda)^{(8)}+(a'_{41})^{(8)}+(p_{41})^{(8)}\right)(q_{40})^{(8)}G^{*}_{40}+(a_{40})^{(8)}(q_{41})^{(8)}G^{*}_{41}\right)$$

$$\left(\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(40)}T^{*}_{41}+(b_{41})^{(8)}s_{(40),(40)}T^{*}_{40}\right)$$

$$\left(\left((\lambda)^{(8)}\right)^{2}+\left((a'_{40})^{(8)}+(a'_{41})^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\right)$$

$$\left(\left((\lambda)^{(8)}\right)^{2}+\left((b'_{40})^{(8)}+(b'_{41})^{(8)}-(r_{40})^{(8)}+(r_{41})^{(8)}\right)(\lambda)^{(8)}\right)$$

$$+\left(\left((\lambda)^{(8)}\right)^{2}+\left((a'_{40})^{(8)}+(a'_{41})^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\right)(q_{42})^{(8)}G_{42}$$

$$+\left((\lambda)^{(8)}+(a'_{40})^{(8)}+(p_{40})^{(8)}\right)\left((a_{42})^{(8)}(q_{41})^{(8)}G^{*}_{41}+(a_{41})^{(8)}(a_{42})^{(8)}(q_{40})^{(8)}G^{*}_{40}\right)$$

$$\left(\left((\lambda)^{(8)}+(b'_{40})^{(8)}-(r_{40})^{(8)}\right)s_{(41),(42)}T^{*}_{41}+(b_{41})^{(8)}s_{(40),(42)}T^{*}_{40}\right)\}=0$$

$$+$$

$$\left((\lambda)^{(9)}+(b'_{46})^{(9)}-(r_{46})^{(9)}\right)\{\left((\lambda)^{(9)}+(a'_{46})^{(9)}+(p_{46})^{(9)}\right)$$

$$\left[\left(\left((\lambda)^{(9)}+(a'_{44})^{(9)}+(p_{44})^{(9)}\right)(q_{45})^{(9)}G^{*}_{45}+(a_{45})^{(9)}(q_{44})^{(9)}G^{*}_{44}\right)\right]$$

$$\left(\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(45)}T^{*}_{45}+(b_{45})^{(9)}s_{(44),(45)}T^{*}_{45}\right)$$

$$+\left(\left((\lambda)^{(9)}+(a'_{45})^{(9)}+(p_{45})^{(9)}\right)(q_{44})^{(9)}G^{*}_{44}+(a_{44})^{(9)}(q_{45})^{(9)}G^{*}_{45}\right)$$

$$\left(\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(44)}T^{*}_{45}+(b_{45})^{(9)}s_{(44),(44)}T^{*}_{44}\right)$$

$$\left(\left((\lambda)^{(9)}\right)^{2}+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\right)$$

$$\left(\left((\lambda)^{(9)}\right)^{2}+\left((b'_{44})^{(9)}+(b'_{45})^{(9)}-(r_{44})^{(9)}+(r_{45})^{(9)}\right)(\lambda)^{(9)}\right)$$

$$+\left(\left((\lambda)^{(9)}\right)^{2}+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\right)(q_{46})^{(9)}G_{46}$$

$$+\left((\lambda)^{(9)}+(a'_{44})^{(9)}+(p_{44})^{(9)}\right)\left((a_{46})^{(9)}(q_{45})^{(9)}G^{*}_{45}+(a_{45})^{(9)}(a_{46})^{(9)}(q_{44})^{(9)}G^{*}_{44}\right)$$

$\left(\left((\lambda)^{(9)} + (b'_{44})^{(9)} - (r_{44})^{(9)}\right)s_{(45),(46)}T^*_{45} + (b_{45})^{(9)}s_{(44),(46)}T^*_{44}\right)\} = 0$

**And as one sees, all the coefficients are positive. It follows that all the roots have negative real part, and this proves the theorem.**

Note: **Same dovetailing explanation holds good for d/dt (partial differential with respect to t), d/dt,d^2/dt^2 (acceralation: double dot). Such an exposition is helpful in optimisation problems, duality of motion, wave motion a homogeneous functions studies, and control theory which we intend to incorporate in future.** It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## Persistence, Coexistence And Thrusty Permanence For Interacting Structured Particles, Transactions In Finance, Econometric Productivity Functions, Environment, Interfaces, Cosmological Interactions Et Al

## Josef Hofbauer And Sebastian J. Schreiber Model Thereof: A Functional Appraisal

ABSTRACT. **JOSEF HOFBAUER AND SEBASTIAN J. SCHREIBER** studied the dynamics of interacting structured systems (in their case population dynamics) can be modeled by $\frac{dxi}{dt} = A_i\ (x)\ xi$ where $x_i\ \epsilon\ R^n_{i,}\ x = (x_1, \ldots, x_k)$, and $A_i(x)$ are matrices with non-negative off- diagonal entries. These models are permanent if there exists a positive global attractor and are robustly permanent if they remain permanent following perturbations of $A_i\ (x)$. Necessary and sufficient conditions for robust permanence are derived using dominant Lyapunov exponents $X_i\ (\mu)$ of the $A_i\ (x)$ with respect to invariant measures $\mu$. The necessary conditions requires $\max_i\ X_i(\mu)>0$ for all ergodic measures with support in the boundary of the non-negative cone. The sufficient condition requires that the boundary admits a Morse decomposition such that $\max_i$ Xi$(\mu)>0$ for all invariant measures $\mu$ supported by a component of the Morse decomposition. When the Morse components are Axiom A, uniquely ergodic, or support all but one population, the necessary and sufficient conditions are equivalent. Applications to spatial ecology, epidemiology, and gene networks are given. Appeared in *Journal of Differential Equations*, 248, 1955-1971 (2010) It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model

becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## INTRODUCTION

A fundamental issue in population biology is what are the minimal conditions to ensure the long-term survivorship for interacting populations whether they be viral particles, bio-chemicals, plants, or animals. When these conditions are met the interacting populations are said to persist or coexist. Since the pioneering work of Lotka and Volterra on competitive and predator prey interactions, Thompson, Nicholson, and Bailey on host-parasite interactions, and Kermack and McKendrick on disease outbreaks, nonlinear difference and differential equations have been used to understand conditions for population persistence [1-5]. One particularly important form of persistence for deterministic models is *permanence* or *uniform persistence* which corresponds to the existence of a global attractor bounded away from extinction of one or more species. When such an attractor exists, interacting populations are able to recover from "vigorous shake ups" of the of the population state. Permanence has been characterized from a topological perspective [6-8] and with average Lyapunov functions [9-12] for differential equation models of the form

$$\frac{dxi}{dt} = x_i f_i(x) \qquad i = 1, \ldots, k \tag{7}$$

Where $x_i$ is the abundance of population $i$, $x = (x_1, \ldots, x_k)$, and $f_i(x)$ is the per-capita growth of population $i$.

An important extension of the concept of the permanence is *robust permanence* [13-16] which requires that (1) remain permanent after sufficiently small perturbations of the per-capita growth rates $f_i$. The importance of robust permanence as one crude form of structural stability for population models are approximations to reality. Consequently, if nearby models (e.g. more realistic models) are not permanent despite the focal model being permanent, and then one can draw few conclusions about the persistence of the biological system being approximated by the model. One can view robust permanence as one crude form of structural stability for debits-

credits particles, propositions, transactions in the celestial subterranean realm and ceratoid dualism, finance like debiting say a current account and crediting to another account, or recrediting to the account which had been earlier debited, further conduction of the cash credit account, black hole formation, guzzling of inventories for the output, creation of income from output, reinvestment of income to investment,gormadizing of matter by black holes, and emission of radiation therefore, gamma ray bursts, in fact most of the transactions and interactions in the universe are conservative and quantum mechanical universe, something like a account acting as predator and prey. More often than not as we have seen from Haimovici model such models could be applied to the problems, propositions, which are not subsumed by the concomitant problems in the bearing. the population models. For this perspective, it is not unexpected that there are permanent systems that cannot be approximated by robustly permanent systems [17]. In [14, 16], criteria for robust permanence were developed with respect to the average per-capita growth rates $\int f_i(x)d\mu(x)$ with respect to invariant probability measures $\mu$ supported on the boundary of the positive orthant (i.e. where one or more Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates are extinct). Roughly, these criteria for robust permanence require that the average per-capita growth rate (i.e. invasion rate) is positive for some missing species.

While (1) can account for many types of Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates. Interactions, it assumes that all individual within populations are exactly the same. However, theoretical biologists have long recognized that different individuals within a population may be in different states (e.g. different sizes or ages, living in different spatial locations) and these differences can have important consequences for population dynamics [18-24]. To account for how these differences influence persistence of interacting populations, we develop criteria for robust permanence for the dynamics of $k$ interacting structured populations. Roughly, these dynamics correspond to replacing $x_i$ in (1) by vectors and $f_i(x)$ by matrices. In section 2, we describe these models in greater detail and review some basic concepts from dynamical systems theory. In section 3, we introduce the structured analog of average per-capita growth rates and define the concept of robustly unsaturated invariant sets. Necessary and sufficient conditions for an invariant set to be robustly unsaturated are proven. For Axiom A or uniquely ergodic invariant sets, the necessary and sufficient criteria are shown to be equivalent. In section 4, we use Morse decompositions and the criteria for robustly unsaturated invariant sets to develop necessary and sufficient conditions

for robust permanence. In sections 5 and 6, we provide several applications of our results to spatially structured ecological models, structured epidemiological models, and models of gene networks. We apply the model to econometrics. To give an example investment needs inventory to be used towards the consummation of output, part of output is to be sold for consolidation of income, and some income might be replenished in the investment and so on. In finance a cash credit account could be debited or credited when there is say purchase of demand bills or discount of usance bills. Cash credit account is debited whenever cheques are issued towards the purchase of raw materials and so on. In fact some of the models have been applied to Organic psychiatry and we have shown how such models are useful in dealing with the propositions of Husserl, Derrida, Deleuze and others.

## 7. MODELS AND ASSUMPTIONS

Let $x_i$ denote the state of the $i$-th population that lies in the non-negative cone $C_i$ of $R^{ni.}$ Define $x=(x_1,...x_k)\epsilon C$ to be the non-negative cone of $R^n$ where $n=\sum_{i=1}^{k} n_i$. Let $C_i^+$ denote positive cone { $x_i \epsilon C_i$: $\prod_i$i noitalupop rof$^j$

$i$ *and* $C^+=\prod_i C^+_i$ , the positive cone for the interacting populations. If $x \mapsto A_i(x)$ is a map into $n_i x n_i$ matrices that describes the growth of population I, then the dynamics of the interacting populations are given by

$$(2)\frac{dx_i}{dt} = A_i(x)x_i \; i=1,...,k$$

Let $x.t$ denote the solution of (2), with the initial condition x, and more generally the semi flow generated by (2). Before stating our assumption on (2), recall a few definitions from dynamical systems. Given sets $I \subseteq R$ and $K \subseteq C$, let $K.I = \{x.t : t \epsilon I, x \epsilon K\}$. A set $K \subseteq C$ *is invariant* if K.t = K for all t>0, and *forward invariant* if $K.t \subset K$ for all $t > 0$. The *omega limit* set of a set $K \subseteq C$ equals $\omega(K) = \bigcap_{t\geq 0,} \overline{K.[t,\infty)}$. The *alpha limit* set of $K \subseteq C$ is $a(K) = \bigcap_{t\leq 0} \overline{K.(-\infty,t]}$. Given a forward invariant set K, $B \subset K$ is an attractor for the semi flow $x.t$ *restricted* to K provided there exists an open neighborhood $U \subseteq K$ of B such that $\omega(U) = B$. The stable set of a *compact invariant* set $K$ is defined by

$$W^s(K) = \{x \epsilon C : \omega(x) \neq \emptyset \text{ and } \omega(x) \subseteq K\}.$$

The semi flow generated by (2) is *dissipative* if there exists an attractor B with $W^s(B) = C$.

Throughout this paper, we make the following assumptions:

**A1**: $x \mapsto A_i(x)$ is continuous

**A2:** $A_i(x)$ is irreducible and has non-negative off- diagonal entries,

**A3:** $x.t$ is defined for all $t \geq 0$

**A4:** $x.t$ is dissipative with global attractor $\Gamma(A)$.

A1 is a basic regularity assumption. The non-negative of off- diagonal entries in A2 implies that there are no negative feedbacks between individuals of different states in population *i*. This assumption is meet for many types of structured models as discussed in sections 5 and 6. The irreducibility assumption A2 is generally meet and implies that all individuals within a population can pass through all states. In Remark1, we discuss how this irreducibility assumption can be relaxed. A3 ensures that the population dynamics are defined for all future time. A4 requires that population densities/ abundances eventually are uniformly bounded, a condition that should be meet for an biologically realistic model.

We say that the semi flow of (2) is *permanent* if this semi flow is dissipative and there is a positive attractor $B \subset C^+$ such that $W^s(B) \subset C^+$. We will denote this positive attractor by $\Lambda(A)$. Permanence ensures that populations can recover from rare large perturbations and allows for a diversity of dynamical behaviors.

For $K \subset C$ and $\delta > 0$, define the $\delta$- *neighborhood* of K as

$$N\delta(K) = \{x \in C : \backslash x - y < \delta \text{ for some } y \in K\}.$$

We define a $\delta$- *perturbation of (2)* to be a system of the form

$$\frac{dx}{dt} = \tilde{A}(x)x$$

that satisfies assumptions A1-A4, $\Gamma(\tilde{A}) \subset N(\Gamma(A))$, *and* $||A(x) - \tilde{A}(x)|| \leq \delta$ *for all* $x \in N(\Gamma(A))$. (2) is *robustly permanent* if there exist >0 and $\epsilon > 0$ such that all $\delta$- perturbations of (2) are permanent, and $d(\Lambda(\tilde{A}), C\backslash C^+) > \epsilon$ for all $\delta$- perturbations i.e., there is a common/ uniform region of repulsion around the boundary.

## INUNDATION OF COMPACT SETS

We begin by studying the linear skew product flows on $\Gamma(A) \times R^n_i$ defined by $(x,y).t = (x.t, B_i(t,x)y)$ where $Y(t) = B_i(t,x)$ is the solution to $\frac{dY}{dt}(t) = A_i(x.t)Y(t)$ with y(0) equal to the identify matrix. Our assumption that $A_i$ is irreducible with non- negative off diagonal entries implies that $B_i(t,x)C_i \subset C_i^+$ for all $x$ and $t > 0$ (see, e.g., [25] ). A

result of Ruelle [26, Prop, 3.2] implies that there exist continuous maps $u_i\, v_i : \Gamma(A) \rightarrow C_i^+$ with $u_i(x)|=|v_i(x)|=1$ such that The line bundle $E_i(x)$ spanned by $u_i(x)$ is invariant i.e. $E_i(x.t)= B_i(x,t)E_i(x)$ for all t ≥ 0.The vector bundle $F_i(x)$ perpendicular to $v_i(x)$ is invariant i.e. $F_i(x.t)= B_i(x,t)F_i(x)$ for all $t \geq 0$.There exist constants $a> 0$ and $\beta>0$ such that(3) $||B_i(t,x)F_i(x)|| \leq a\, exp(-\beta t)||B_i(t,x)E_i(x)||$ For all $x \in K$ and $t \geq 0$.**Interaction rates.** Given any x $\epsilon$ $\Gamma$(A), define the *interaction and inundation rate* of *species i at particle/proposition/problem/ state x as* $\lambda_i(x)= limsup_{t\rightarrow\infty}\left(\frac{1}{t}\right)$ *in* $||B_i(t,x)||$ when $x_i$ =0, $\lambda_i(x)$ provides an upper bound to the rate of growth of population *i* when introduced at infinitesimally small densities. Important properties of this invasion rate are summarized in the following two propositions. For instance, the second proposition implies if $(x.t)_i$ stays bounded away from zero (i.e. population *i* persist), then

$\lambda_i(x)=0$. Alternatively, if $\lambda_i(x)<0$,then population *i* is doomed to extinction. **Proposition 1.** $\lambda_i(x)$ *satisfies the following properties:* $\lambda_i(x)= limsup_{t\rightarrow\infty}\left(\frac{1}{t}\right) In \quad |B_i \quad (t,x)v|$ *for any* $v>0$ *in* $R^{ni}$, *and* $\lambda_i(x)= limsup_{t\rightarrow\infty}\left(\frac{1}{t}\right) \int^t_0 (A_i(x.s)u_i(x.s),u_i(x.s))ds.$

**Proof:** To prove the first property, we first show that

$$\lambda i(x) = limsup_{t\rightarrow\infty} \frac{1}{t} In\ |B_i(t,x)u_i(x)|.$$

To this end, let $v \in R^{ni}$ be any non-zero vector. Since $R^{ni}= E_{i(x)} \quad F_i(x)$, we can write $v=au_i(x)+w$ with $a \in R$ and $w \in F_i(x)$. Equation (3) implies

$$|B_i(t,x)v| \leq a|\, B_i(t,x)u_i(x)|+ |B_i(t,x)w| \leq |\, B_i(t,x)u_i(x)|(a+ \alpha+ exp\ (-\beta t)|\, w|).$$ Hence, $limsup_{t\rightarrow\infty} \frac{1}{t} In\ B_i(t,x)\ v| \leq$

$$limsup_{t\rightarrow\infty} \frac{1}{t} In\ B_i\,(t,x)\ u_i\,(x) \leq \lambda_i(x)$$

for all non zero vectors $v \in R^{ni}$. Since $||B_i(t,x)|| = sup|v|=1B_i(t,x)v|$, this inequality implies that $\lambda i(x) = \limsup_{t\rightarrow\infty} \frac{1}{t}$ In $|B_i(t,x)v|$, this inequality implies that $\lambda_i(x) = limsup_{t\rightarrow\infty}\left(\frac{1}{t}\right)$ In $|t,x)u_i(x)|$. Now let v $\epsilon$ $C_i^+$. Then, we can write $v = au_i(x)+ w$ with $a$ >0 and $w \in F_i$.

Equation (3) implies

$$|B_i(t,x)v| \geq a|B_i(t,x)u_i(x)|-|B_i(t,x)w|$$

$$\geq |\, B_i(t,x)u_i(x)|(a- \alpha+ exp\ (-\beta t)|\, w|).$$

Since a>0, this inequality implies that

$$\lambda i(x) \geq \limsup_{t\to\infty} \frac{1}{t} \text{In } |B_i(t,x)v| \geq \lambda i(x) \geq \limsup_{t\to\infty} \frac{1}{t} \text{In } |B_i(t,x)u_i(x)| = \lambda i(x)$$

To prove the second assertion, let $b_i(t,x) = In\ |B_i(t,x)u_i(x)|$. Invariance of $u_i(x)$ implies that $b_i(t,x)$ is additive:

$$B_i(t+s,x) = In|B_i(t+s,x)\ u_i(x)| = In\ |B_i(s,x.t)B_i(t,x)u_i(x)$$

$$= \text{In } |B_i(s,x.t)\ u_i(x.t)||\ B_i(t,x)\ u_i(x)|$$

$$= \text{In } b_i(s,t.x)\ b_i(t,x)$$

Additivity of $b_i(t,x)$ and the fact that $b_i(0,x)$ =0 implies

$$\frac{d}{dt} b_i(t,x) = \lim_{s\to 0} \frac{bi(t+s,x-bi(t,x)}{S-}$$

$$= \lim_{s\to 0} \frac{bi(s,x.t)}{S} = \frac{d}{ds}|_{s=0}\ b_i(s,x.t)$$

=

$$=(A_i(x.t)u_i(x.s),u_i(x.s),$$

The Fundamental Theorem of Calculus implies

$$B_i(t,x)= \int_0^t (A_i(x.s)\ u_i(x.s),\ u_i(x.s))ds\backslash$$

and the second assertion follows.

**Proposition 2:** *For the solutions of (2) we have:* if $x_i>0$, then $\lambda i(x) \leq 0$. If $\lambda i(x)<0$, then $\lim_{t\to\infty} (x.t)_i=0$ If $limsup_{t\to\infty}$ $(x.t)_i|>0$, then $\lambda i(x)=0$.

**Proof:** First, assume that x satisfies $x_i>0$. Since the semi flow is dissipative, there exists $\gamma >0$ such that $|x.t| \leq \gamma$ for all $t \geq 0$. Proposition 1 and the definition of the skew product flow imply

$$\lambda i(x) = \limsup_{t\to\infty} \frac{1}{t} \text{In } |B_i(t,x)x_i| = \limsup_{t\to\infty} \frac{1}{t} \text{In } |\ (x.t)_i\ |$$

$$\leq \limsup_{t\to\infty} \frac{In\ \gamma}{t} = 0$$

To prove the second assertion, assume that $\lambda i(x) < 0$. *If* $x_i$=0, then the invariance of the faces of C imply that $(x.t)_i$=0 for all t≥ 0. Alternatively, if $x_i$> 0, then Proposition I and the skew product flow imply

$$\limsup_{t\to\infty} \frac{1}{t} In|(x.t)_i| = \lambda i(x) < 0.$$

The final assertion follows from the first two assertions. 5

**3.2. Invariant measures.** We review some definitions from ergodic theory. Given a Borel probability measure $\mu$ on ***C***, the *support* of $\mu$, denoted supp( $\mu$),is the smallest closed set whose complement has measure zero. A Borel probability measure $\mu$ is called *invariant* for (2) provided that Integral $h(x)d\mu(x)$ = Integral $h(x.t)d\,\mu(x)$for all $t \geq 0$ and for all bounded continuous functions $h$:C→R. An invariant measure $\mu$ is called *ergodic* provided that $\mu(B)$=0 or 1 for any invariant Borel set *B.*

For an invariant measure $\mu$ define the invasion rate species $i$ with respect to $\mu$ as

$$\lambda_i(\mu) = \int_c \langle \mathrm{A}_i(x)\mathrm{u}_i(x), \mathrm{u}_i(x)\rangle d\mu(x)$$

**Proposition 3:** Let $\mu$ be an *invariant measure.* $\lambda_i(\mu)$ *Satisfies the following properties:*$\lambda_i$(x) =lim $_{t\to\infty}$ $\frac{1}{t}\int_0^t \langle \mathrm{A}_i(x.s)\mathrm{u}_i(x.s), \mathrm{u}_i(x.s)\rangle, ds$ exists for $\mu$ - *almost every* $x$. *Moreover, if* $\mu$ *is ergodic,* $\lambda_i$(x) = $\lambda_i(\mu)$ for $\mu$- *almost every* $x$.*If* $\mu$ *is ergodic, then there exists* $I \subset \{1, \dots, k\}$such that $\mu$ ( $\Pi_{i\in I} \boldsymbol{C}_i^+$)=1 and $\lambda_i(\mu) = 0$ *for all* $i \in I$.

**Proof:** Let $\mu$ be an invariant measure. Define $\mathrm{h}_i(x) = \langle \mathrm{A}_i(x)\mathrm{u}_i(x), \mathrm{u}_i(x)\rangle$. The Birkhoff's' ergodic theorem and Proposition 1 imply

$$\lambda_i(x) = \lim_{t\to\infty} \frac{1}{t} \int_0^t \mathrm{h}_i(x.s)\mathrm{ds}$$

Exists for $\mu$- almost every $x$. Moreover, $\lambda_i(x) = \lambda_i(\mu)$ $\mu$- almost surely if $\mu$ is ergodic.

Assume $\mu$ is ergodic. By ergodicity and invariance of the faces of **C**, there exists $I \subset \{1, \dots, k\}$ such that $\mu$ ( $\Pi_{i\in I} \boldsymbol{C}_i^+$)=1. Let K $\subset \Pi_{i\in I} \boldsymbol{C}_i^+$ be a compact set such that $\mu(K) > 0$. The Poincare recurrence theorem and

the Birkhoff ergodic theorem imply that there is $x \in \Pi_{i\in I} \boldsymbol{C}_i^+$ such that $x.t_n \in K$ for some $t_n \uparrow \infty$ and $\lambda_i(x) = \lambda_i(\mu)$ for all $i$. The third assertion of Proposition 2 implies $\lambda_i(\mu) = 0$ for all $i \in I$.

**3.3. Robustly unsaturated sets**. Let $K \subset \partial C$ be a compact isolated invariant set for the flow $x.t$ restricted to $\partial C$. $K$ is unsaturated if $W^s(K) \subset \partial C$ and $K$ is isolated for $x.t$. If $K$ is not unsaturated, then $K$ is saturated. $K$ is robustly unsaturated for (2) if there exists $\delta > 0$ such that the continuation of $K$ for any $\delta -$ perturbation of (2) is unsaturated.**Theorem 1:** Let $K$ be a compact isolated invariant set for $x.t$ restricted to $\partial C$. If one of the following equivalent conditions hold For all invariant measures $\mu$ supported by $K$ (4) $\max_{1\le i\le k} \lambda_i(\mu) = 0$

There exist $p_i > 0$ such that (5) $\sum_{1\le i\le k} p_i \lambda_i(\mu) = 0$

For all ergodic probability measures supported by $K$.

then $K$ is robustly unsaturated for (2). Alternatively, if $x \to A(x)$ is twice continuously differentiable and $K$ is robustly unsaturated, then (4) holds for all ergodic measures $\mu$ supported by $K$.

**Remark 1**: For some applications (e.g., the disease model considered in section 6.1) it useful to relax the irreducibility assumption A2. For instance, if there exists an open neighborhood U of K such that for each $i$, $A_i(x)$ has a fixed off-diagonal sign pattern for all $x \in U$,then $A_i(x)$ can be decomposed into a finite number, say $m_i$, of irreducible components. For each of these irreducible components, one can define $\lambda_i^j(\mu) = \int_C \langle A_i^j(x) u_i^j(x), u_i^j(x)\rangle \, d\mu(x)$where $A_i^j(x)$ is the submatrix of $A_i(x)$ corresponding to the $j$-th irreducible component of $A_i(x)$ $and$ $u_i^j(x)$ is the continuous invariant sub bundle for the irreducible linear cocycle determined by $A_i^j(x)$. If we define $\lambda_i(\mu) = \max_{1\le j\le m_i} \lambda_i^j(\mu)$, then all of the assertions of Theorem 1 still hold.

**Proof:** To prove the first assertion of the theorem, we need the following lemma. We call an invariant measure $\mu$ for (2) saturated if $\lambda_i(\mu) \le 0$ for all $i$

**Lemma 1**: Let $\delta_n$ be a non –negative sequence that converges to zero as n$\to\infty$. If $\mu_n$ are saturated invariant measures for $\delta_n -$ perturbations of (2), then the weak$^*$ limit points of $\{\mu_n\}_{n=1}^{\infty}$ is a non-empty set consisting of saturated invariant measures for (2).

**Proof:** Let $\delta_n$ be a non-negative sequence that coverages to zero. Let $\mu_n$ be saturated invariant measures for $\delta_n$ perturbations, $\dot{x} = \mathrm{A}^n(x)x$, of $\dot{x} = \mathrm{A}(x)x$. By weak$^*$ compactness of Borel probability measures supported on $N_1(\Gamma(\mathrm{A}))$, there exist weak$^*$ limit points of $\mu_n$. Let $\mu$ be such a weak$^*$ limit point. Since $\mu_n$ are supported by $N_1(\Gamma(\mathrm{A}))$, for all $n$ sufficiently large, $\mu$ is supported by $N_1(\Gamma(\mathrm{A}))$.To verify that $\mu$ is an invariant measure for (2),let $h: C \to \mathrm{R}$ be a bounded continuous function. Let $t > 0$ and $\in > 0$ be given. Choose $n$ sufficiently large so that $|\int_c \mathrm{h}(x)\, d\mu(x) - \int_c \mathrm{h}(x) d\mu_n(x)| \leq \in$, $|\int_c \mathrm{h}(x.t)\, d\mu(x) - \int_c \mathrm{h}(x.t) d\mu_n(x)| \leq \in$, and $|\mathrm{h}(x.t) - h(x^n.t)| \leq \in$ for all $x \in N_1(\Gamma(\mathrm{A}))$.Then $|\int_c \mathrm{h}(x.t) - \mathrm{h}(x) d\mu(x)|$ is

$\leq \quad |\int_c \mathrm{h}(x.t)\, d\mu(x) - \int_c \mathrm{h}(x.t)\, d\mu_n(x)| + |\int_c \mathrm{h}(x.t)\, d\mu_n(x) - \int_c \mathrm{h}(x)\, d\mu(x)|$

$\leq \quad \in + \int_c |\mathrm{h}(x.t) - h(x^n.t)|\, d\mu_n(x)| + |\int_c h(x^n.t)|\, d\mu_n(x) - \int_c \mathrm{h}(x) d\mu(x)|$

$\leq \quad 2 \in + |\int_c h(x^n.t) - h(x)\, d\mu_n(x)| + |\int_c h(x) d\mu_n(x) - \int_c h(x) d\mu(x)|$

$\leq \quad 3 \in$

Where the last line follows from the invariance of $\mu_n$ for $x^n.t$. Since $\in > 0$ and t > 0 are arbitrary, we have $\int_c h(x.t) d\mu(x) = \int_c h(x) d\mu(x)$ for all t >0 and all bounded continuous functions $h: C \to R$. It follows that $\mu$ is an invariant measure for (2). To see that $\mu$ is saturated, define $\mathrm{h}_i^n(x) = \langle \mathrm{A}_i^n(x)\mathrm{u}_i^n(x), \mathrm{u}_i^n(x)\rangle$ where $\mathrm{u}_i^n(x)$ spans the invariant one-dimensional bundle given by $[26, Prop.3.2]$ for $\frac{dx}{dt} = A^n(x)x$. Since $\mathrm{A}_i^n(x) \to \mathrm{A}_i(x)$ uniformly for $x \in N_1((\mathrm{A}))$ as $n \to \infty$, $[26, Prop.3.2]$ implies that $\mathrm{u}_i^n(x) \to \mathrm{u}_i(x)$ converges uniformly for $x \in N_1(\Lambda(\mathrm{A}))$. Given $\in > 0$, choose $n$ sufficiently large so that $|\mathrm{h}_i^n(x) - \mathrm{h}_i(x)| \leq \in$ for all $x \in N_1(\Lambda(\mathrm{A}))$ and $|\int_c \mathrm{h}_i^n(x)\, d\mu(x) - \int_c \mathrm{h}_i^n(x)\, d\mu_n(x)| \leq \in$.

Then

$$\lambda_i(\mu) \leq \int_c |\mathrm{h}_i(x) - \mathrm{h}_i^n(x)| d\mu(x) + |\int_c \mathrm{h}_i^n(x)\, d\mu(x) - \int_c \mathrm{h}_i^n(x)\, d\mu_n(x)|$$

$+ \int_c h_i^n(x)\, d\mu_n(x)$

$\leq 2\epsilon$

Since $\epsilon > 0$ is arbitrary, $\lambda_i(\mu) \leq 0$ for all $i$ and $\mu$ is saturated.

Using this lemma, we prove that if K is saturated, then there exists a saturated invariant measure $\mu$ supported by K. Assume K is saturated. Work of Hofbaue and So $[7, Thm.\ 2.1]$ implies that either there exists $\mathcal{Y} \in W^s(K) \cap C^+$ or K is not isolated for the unrestricted flow $x.t$ If there exists $\mathcal{Y} \in W^s(K) \cap C^+$, then for all $t \leq 0$ define $v_t = \frac{1}{t} \int_0^t \delta_{y.s}\, ds$ where $\delta_{y.s}$ denotes a Dirac measure based at the point $\mathcal{Y}.s$. Dissipativeness of (2) and weak* compactness of the Borel probability measures supported by (A) imply there exist $t_k \to \infty$ such that the sequence $v_{tk}$ converges in the weak* topology to a Borel probability measure $\mu$ with support in K.A standard argument implies that $\mu$ is $x.t$ invariant. Define $h_i: C \to R$ by $h_i(x) = (A_i(x)u_i(x), u_i(x))$.Proposition 2 and weak* convergence imply that

$$0 \geq (\lambda i(y) = \limsup_{t\to\infty} \frac{I}{t} \int_0^t h_i(y.s)ds$$

$$= \limsup_{t\to\infty} \int_c h_i\, dv_t$$

$$\geq \limsup_{t\to\infty} \int_p h_i\, d\, v_{tk} = \lambda i(\mu)$$

For all $i$. Alternatively suppose that K is not isolated for the semi flow Then there exists a sequence of positive $\omega - limit$ sets that accumulate on K. Let $\mu_n$ be a sequence of ergodic probability measures supported by these $\omega -$ limit sets. Propositions 1 and 3 imply that $\mu_n$ are saturated for all $n$ Applying Lemma I with $\delta_n$=0 for all $n$ implies that there exists a saturated invariant measure $\mu$ suppoorted by *K*. Hence we have shown that if *K* is saturated then there exists a saturated invariant probability measure $\mu$ supported by *K* Equivalently, (4) holding for all invariant probability measures with support in *K* implies that *K* is unsaturated.

Next we show that if *K* not robustly saturated, then there exists a saturated invariant measure supported by *K* Indeed suppose *K* is not robustly saturated, Then there exists a non-negative sequence $\delta_n$ converging to zero and a

sequence of saturated measures $\mu_n$ for $\delta_n$-perturbations of (2)with support in the continuation of *K*. Let $\mu$ be a weak* limit point of $\{\mu_n\}\infty_{n=1}$ Lemma 1 implies that $\mu$ is saturated. Moreover, since the continuation of *K* converges to *K* as $\delta_n \to 0, \mu$ is a saturated invariant measure for (2) supported by *K* Hence we have shown that if *K* is not robustly saturated then there exists a saturated invariant measure supported by *K* Equivalently if (4) holds for all invariant measures with support in *K* then *K* is robustly unsaturated.

To see the equivalence of the conditions given by (4) and (5) let $\Delta = \{p \in R_+^k : \sum_i P_i = 1\}$ and notice that

$$\text{Min max } \lambda i\ (\mu = \min_{\mu}\ \max_{p\epsilon\Delta}\ \sum_i\ Pi\lambda i(\mu)$$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K* The Minimax theorem (see, e. g.,[27] implies that

$$(6)\ \min \max\ \lambda i(\mu) = \min_{\mu}\ \max_{p\epsilon\Delta}\ \sum_i\ Pi\lambda i(\mu)$$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K.* Since min$\mu$ $\sum_i\ P_i\ \lambda_i(\mu)$ is attained at an ergodic probability measure with support in *K*, the equivalence of the conditions given by (4) and (5) is established.

To prove the final assertion of the Theorem assume $x \to A(x)$ is twice continuously differentiable and there exists a saturated ergodic measure $\mu$ supported by *K* Proposition 3 implies that there exists *I*C {1,…,k} such that (□$i \in iC_i^+$)=1and $\lambda i(\mu) = 0$ For all $i \in I. SinceKc\ .\partial$□,$\{1, \ldots . k\}\backslash I$ is non-empty. We will show that (2) is not robustly permanent by proving that for all $\delta > 0$ there exists a $\delta$ −perturbation of (2) for which the continuation of *K* is saturated Let $\delta > 0$ be given Choose V $c$ N$\delta$ ($\Gamma$(A)) to be a compact neighborhood of Γ(A) such that V.t $c$ int *V* for all t > 0.Let $\eta > 0$ be such that *W*=:*N* $\eta$(Γ(A))$c$ *V* Let p:c→ [0,1]$be$ a$C^\infty$function such that $p(x) = 1 \quad for\ all\ x \in \Gamma(A) and\ p(x) = 0 for\ all\ x \in C\backslash W\ Define\ \tilde{A} = (\tilde{A}_1, \ldots, \tilde{A}_k)$ by

$$\tilde{A}_i(x) = \{\frac{A(x)}{A(x)-\frac{\delta}{2Idp\,(x)}} \quad \frac{if\ i\epsilon I}{if\ i \notin I}$$

Where Id denotes the identity matrix of appropriate dimension Let x. t denote the semi flow of $x = \tilde{A}(x)x$.Since$x.t = x.t$ $whenever\ x.[0,t] \in C\backslash W, it\ follows\ that\ \ \tilde{V}.t \subset intV\ for\ all\ t >$ $0. Hence, \Gamma(\tilde{A}) c\ V\ c\ N\delta\ (\ (A))). We\ \ also\ have\ ||A(x) - \tilde{A}(x)|| \leq \frac{\delta}{2}$Therefore $x = \tilde{A}(x) is\ a\ \delta -$ perturbation of (2) By construcation,$x.t = x.t\ for\ all\ x \in K and\ t \geq 0. Hence, \mu$ is ergodic for the semiflow of $\dot{x} = \tilde{A}(x)x\ and\ \lambda_i(\mu) \leq -\frac{\delta}{2}$ for this semiflow and $i \in \{1, \dots, k\}\backslash I$.Let $\mathcal{L}$ and O be the Lyapunov exponents and Oseledec regular points supported by $\mu$ for $x = \tilde{A}(x)x(see, e.g.,$ [28] for definitions).At each point $x \in O$, the splitting of $R^n$ determines three subspaces: the stable subspace $E^s(x)$, the center subspace $E^c(x)$, and the unstable subspace $E^u(x)$ Proposition 3 and our choice of Ã imply that $E^s(x) \cap C^+ \neq \emptyset$. The Pesin stable manifold theorem [28, Corollaries 3.17 and 3.18] implies that tangent to $E^s(x)$ $E^c(x)$ and $E^u(x)$ are locally $x.t-$invariant families of $C^1$ discs $W^s_x, W^c_x$ and $W^u_x$ corresponding to the stable ,center and unstable manifolds The family of stable manifolds $W^s_x$ is contained in $W^s$ (Ã,supp $(\mu)$). $Since\ E^s(x) \cap C^+ \neq \emptyset$, $W^s_x \cap C^+ \neq \emptyset$ for some $x \in K$. Consequently $K$ is saturated for $x = \tilde{A}(x)x$.

The difference between the sufficient and necessary condition is that for the sufficient condition (4) has to hold for all invariant measures supported by *K*, while for the necessary condition (4) has to hold only for ergodic measures supported by *K*. since the invariant measures lie in the convex hull of the ergodic measures the sufficient condition can be more restrictive then the necessary condition.

For three Classes of invariant sets the necessary and sufficient conditions coincide Recall a compact invariant set *K* for (2) is uniquely ergodic if K only supports one invariant measure Recall a compact invariant set *K* is Axiom A if the flow of (2) restricted to *K* is transitive and hyperbolic (see, e.g.,[29] for definitions.)

**Corollary 1:** Assume $x \to A(x)$ is twice continuously differentiable It K$\subset \partial$□ is a compact invariant set and either is uniquely ergodic or Axiom A, then *K* is robustly unsaturated if and only if (4) holds for all ergodic, measures $\mu$ with support in *K*

Proof. If K is uniquely ergodic then the assertion follows immediately For an Axiom A invariant set Sigmund [30, thm.1] has proven that invariant measures supported by periodic orbits of $K$ are dense in the set of invariant probability measures supported by $K$ In particular any invariant measure can be approximated by an ergodic measure and consequently the result follows.

Another special case where the necessary and sufficient conditions coincide is when $K$ supports all populations except one for discrete time models this case was considered by [22].

**Corollary2:** Assume $x \to A(x)$is twice continuously differentiable Let K$c\ II_i \geq 2C_i^+$ be a compact invariant set. Then K is robustly unsaturated if and only if $\lambda_1(\mu) > 0$ for all invariant measures supported by $K$

## 4. MORES DECOMPOSITIONS AND ROBUST PERMANENCE

To state the sufficient condition for robust permanence, we use a characterization of permanence due to Garay [8] and Hofbauer and So [7] that involves Morse decompositions of the boundary flow Conley[31] defined a collection of sets$\{M_{1,\ldots,}M_m\}$ to be a Morse decomposition for a compact invariant set $K$ if $M_1\ldots,M_m$ are pairwise disjoint, compact isolated invariant sent for the flow of (2) restricted to K. For each $x \in K$ there are integers r=r$(x) \leq s = s(x)$ such that $a(x) \subseteq M_r$and $\omega(x) \subseteq M_{s.}$ If $r(x) = s(x), then\ x \in M_{r.}$ Garay, Hofbauer and so [7,8] proved the following characterization of permamnce

**Theorem 2** (Garay, Hofbauer-so) if $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$?then (2) is permanent if and only if each of the components $M_1$ are unsaturated.

Theorem 1 and 2 imply the following result: As stated in the above all these results are true for most of problems, interactions, transactions in cosmology, econometrics, finance and other subjects, which we shall illustrate as we go on constructing models for the same.

**Theorem 3:** If $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$?and (4) holds for each of the components of the More decomposition, then (2) is robustly permanent Conversely if $x \to A(x) is\ twice\ continuously\ differentiable$

*and (4) is violated by an ergodic measure supported by one of the components of the Morse decomposition, then (2) is not robustly permanent.* Theorem 3 in conjunction with Corollaries 1 and 2 yield a characterization of robust permanence for a class of structured models.

**Corollary 3:** *Let $\{M_1,.....,M_m\}$ os a Morse decomposition for $\partial C$. Assume $x \mapsto A(x)$ is twise continuously differentiable and for each Morse component $M_i$ one of the following assertions hold $M_i$ is Axiom A, $M_i$ is uniquely ergodic, or There exists $j \in \{1,....,k\}$ such that $M_i$ ----- $\{ x \in C : x_l \gg 0$ for $l \neq j\}$ Then (2) is robustly permanent if and only if (4) holds for all ergodic measures supported by $U_i M_i$.*

Our result also provided a structured analogue for characterizing totally permanent systems[32]. For an ergodic probability measure$\mu$, define sup ($\mu$) to be the subset $I$ -----$\{1,\ldots,k\}$ such that $\mu(---_{i \in I} C_+) = 1$

**Corollary 4:** *The following statements are equivalent: For all ergodic probability measures $\mu$ with support in $\partial C$, $\lambda_i(\mu) > 0$ for all $I \in \{1,...,k\} \setminus supp(\mu)$ (2) and all of its subsystems are robustly permanent.*

*Proof. Suppose the first statement holds. Then unsaturated condition given by (5) with p= (1,...,1) and Theorem 3 implies that (2) and all its subsystems are robustly permanent.. The other direction follows form the necessary condition for robust permanence in Theorem 3 and the second assertion of Proposition 3.*

*To illustrate the broad applicability of theorem 3, we develop applications to spatially structured models in section 5 and to a disease model and a gene network model in section 6.*

## 5. APPLICATIONS: PATCH MODELS

A fundamental application of our result are to spatially structured models, with *k* species dispersing between *m* patches {33-35}:

(7) $$\frac{dx_i^j}{dt} = x_i^j f_i^j(x^j) + \sum d_i^{jl} x_i^l - e_i^j x_i^j \qquad i = 1,..,k \qquad j=1,\ldots,m$$

*l*

Where $x_i^j$ denotes the density of species I in patch $j$, $x^j = (x^j_1, \ldots, x^j_k)$ is the vector of species densities in patch $j$, $f^j_i$ is the per-capita growth rate of species $i$ in patch $j$, $d_i^{jl} \geq 0$ is the dispersal rate for species $i$ from patch $l$ into patch $j$, and $e_i^j$ is the emigration rate of species $i$ out of patch $j$. Hence $e_i^j \geq \sum_l d_i^{lj}$. We assume that the matrices $(d_i^{lj})_{j,l}$ are irreducible for each species $i$. it is then easy to write (7) in the form (2), with $x_i = (x^l_i, \ldots, x_i^m)$ 'where' denotes transpose.

For a single species (7) generates a monotone flow. Under mild assumption, e.g., each $f^j_i$ *is* decreasing and negative for large densities, there is a globally stable equilibrium, see, e.g.,[36] and [37, sect. 5.4]. This equilibrium will be positive if the invasion rate at the origin $\lambda(0)$, which is given by the leading eigenvalue (stability modulus) of the matrix $(d^{jl} + f^j(0) - e^j)$, is positive.

5.1 Two types of debitable and creditable systems:. For *two competing species,* e.g., $f^j_i$ is decreasing with respect to both species and is negative at large densities, the dynamics are still monotone, and hence almost all orbits converge to an equilibrium. Robust permanence requires two conditions. First, $\lambda_i(0) > 0$ for both species in which case there are two single species equilibria $E_1$ and $E_2$. Second, the invasion rates $\lambda_2(E_1)$ and $\lambda_2(E_2)$ are positive. Whether there is a unique positive globally stable equilibrium for the robustly permanent system depends in a delicate way on the system parameters, see [38]. For slowly dispersing populations (i.e. $0 < e^j_i \ll 1$ for all $i, j$), the robust permanence condition is particularly

Straight forward to verify. Let $E^j_i$ be the largest solution to $x^j_i \quad f_i^j \ (x^j_i) = 0$ i.e. the equilibrium attained by species $I$ in patch $j$ when there is no dispersal and no competitors. For slowly dispersing populations, $E_1$ is close to the product $(E_1^1, \ldots, E_1^m)$, and the matrix $A^2(E_1)$ is close to the diagonal matrix $\mathrm{diag}(f^i_2(E_1{}^j))^m_{j=1.}$ Hence the invasion rate $\lambda_2(E_1)$ is close to $\max_j f_2^j(E^j_1) > 0$. So we obtain: The system with two slowly dispersing competing species is robustly permanent if and only if system refers to system of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

(8) $\max_j f_2^j(E^j_1) > 0$ and $\max_j f_2^j(E^j_2) > 0$

Intuitively, robust permanence requires for each species there is at least one patch where it can persist.

For *systems, particle interactions, financial transactions, cosmological interactions and reactions, as we shall show model by model the* situation are even easier. Since (under mild assumption on the prey dynamics) the Morse decomposition of $\partial C$ is given by two equilibria $\{0,E_1\}$, the spatial predator-prey system is robustly permanent, if the predator can invade at the prey equilibrium: $\lambda_2\ (E_1) > 0$. However, the global dynamic are likely to be more complicated. For instance, if in each patch there is a globally stable limit circle (e.g.,[39-41]) and dispersal is sufficiently slow (i.e. $0< e^j{}_i \ll 1$for all *i, j*), our result about robust permanence and the theory of normally hyperbolic manifolds [42] imply that there is a positive *m* dimensional torus, which attracts almost all orbits in C, whenever $\max_j f_2^j (E^j_1)>0$

**5.2 Rock-Paper-Scissors Dynamics.** The Lotka-Volterra model of rock-paper-scissor dynamics is a simple model that is used as property for understanding intransitive ecological outcomes [43,44]. Here, a simple spatial version of this dynamic is given by

$$\frac{dx_1^j}{dt} = x_1^j(1-x_1^j - \beta^j\ x_2^j\ - \alpha^j\ x_3^j\ ) + \sum d_1^{jk}\ \ x_1^k\ - e_1^j\ x_1^j$$

$$\frac{dx_2^j}{dt} = x_2^j(1-\alpha^j\ x_1^j\ - x_2^j - \beta^j\ \ x_3^j\ ) + \sum d_2^{jk}\ \ x_2^k\ - e_2^j\ x_2^j$$

$$\frac{dx_3^j}{dt} = x_i^j(1-\beta^j\ x_1^j\ - \alpha^j x_2^j - \beta^j\ \ x_3^j\ ) + \sum d_3^{jk}\ \ x_3^k\ - e_3^j\ x_3^j$$

Where $a^j \in (0,1)$ and $\beta^j > {}^1$ for all *j*. A more general version of this dynamic is presented in [44]. Under the assumption that there is no cost to dispersal (i.e. $e^j_i = \sum_l d^{l,j}_i$ for all *i, j*) , the maximal invariant set in $\partial C$ consist of the origin 0 and a heteroclinic cycle connecting positive single species equilibria$E_1$, $E_2$ and $E_3$. For slowly dispersing population ($0< e^j_i \ll 1$ for all *i,* j), $E_1$ is close to $(x_1, x_2, x_3) = (1,\ldots,1,0,\ldots,0,0,\ldots,0)$ and the invasion rates $\lambda_2\ (E_1)$and $\lambda_3\ (E_1)$are close to $\max_j 1-a^j$ and $\max_j 1 - \beta^j$. For the other equilibria $E_i$, the debitable rates of the missing species are also given by $\max_j 1-a^j$ and $\max_j 1 - \beta^j$. Consider the Morse decomposition of $\partial C$ given by 0 and the heteroclinic cycle. Since $\lambda i\ (0) = 1> 0$ for $i = 1, 2, 3, 0$ is robustly insaturated. An algebraic computation reveals that (4) holds for all invariant supported by the heteroclinic cycle (i.e. all convex combination of the Dirac measures supported by the equilibria) if and only if

$$\max(1-a^j) > -\max\ (1 - \beta^j\ )$$

$j$ $j$

Equivalently,

$$2 > \min a^j + \min \beta^j$$

$j$ $j$

In particular, even if the heteroclinic cycle is attracting for each patch when the system is uncoupled (i.e. $2 < a^j + \beta^j$ for all $j$ ), it can be repelling for the weakly coupled system.

5.3. Three species **Lotka Volterra in spatially homogenous environments.** Consider Lotka-Volterra dynamics in a spatially homogenous environment: $(x + a)^n =$

(9) $$\frac{dx_i^j}{dt} = x_i^j (r_i - \sum_s a_{is} x_s^j) + \sum_l d_i^{jl} x_i^l - e_i^j x_i^j , \quad i = 1, \dots, k; \; j = 1, \dots, m$$

We assume that dissipative can be shown by a linear Liapunov function: There are $c_i > 0$ such that $\sum_i c_i a_{is} x_i x_s < 0$ holds for all $x \in R_+^k \setminus \{0\}$. Then the weighted sum of densities across all patches $\sum_{i,j} c_i x_i^j$ is decreasing for large densitities and hence (9) generates a dissipative semiflow with global attractor $r$ . We further assume that for each patch *j* and each species *i* the immigration rate equals the emigration rate , i. e,

(10) $$e_i^j = \sum_l d_i^{jl}$$

This guarantees that the set H= $\{x \in C: x_i^j = x_i^l \forall i, j, l\}$ of spatially homogenous states is forward invariant under (9).

Now consider the Lotka-Volterra dynamics without spatial structure

(11) $$\frac{dx_i}{dt} = x_i(r_i - \sum_s a_{is}x_s), \qquad i = 1, \dots, k$$

(11) $$\frac{dx_i}{dt} = x_i(r_i - a_{is}x_s), \quad i = 1, \dots, k$$

It has been shown [12,ch.16.1,16.2] That for $k$ =3,(11) is robustly permanent (robust meaning here within the class of Lotka –Volterra systems) if and only if all equilibria on the boundary of $R^3_+$ are unsaturated and whenever there is a heteroclinic cycle connecting the one species equilibria then this cycle is repelling. Equivalently the following four conditions hold. There exists an interior equilibrium $\dot{x}$ (i.e $A\dot{x} = r \;\; with \; \dot{x} \gg 0$); Det (-A)>0; the 2 species subsystems are not bistable competition systems if there is a heteroclinic cycle between the one species equilibria, say $\mathbf{E_1 \to E_2 \to E_3 \to E_1}$ then the following inequality holds

12) $$\lambda_2(E_1)\lambda_3(E_2)\lambda_1(E_3) > |\lambda_3(E_1)\lambda_1(E_2)\lambda_2(E_3)|$$

Furthermore, the proof in [12, ch.16.1, 16.2] shows that the boundary flow of (11) has a simple Morse decomposition: if (iv) applies then there are two Morse sets, the origin 0 is a repeller, and the heteroclinic cycle is the dual attractor (within$\partial R^3_+$); if (iv) does not apply then the finitely many boundary equilibria form a Morse decomposition of the boundary flow.

**Species, population we mean types of and population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.**

**Theorem 4:** Under the above assumptions,(9) for k=3 is permanent if and only if *(i),(ii),(iii)* and *(iv)* hold.

Proof: Suppose (9) is robustly permanent. Since the nonspatial system (11) is the restriction of (9) to the invariant subspace **H**, (11) must be robustly permanent and hence*(i)-(iv)* hold

Conversely, suppose that (i)-(iv) hold. We show that the above Morse decomposition of the boundary flow of (11) (now in **H**) is also Morse decomposition for (9) restricted to the maximal invariant subset of $\partial C$ which is contained in the spatial two species subsystems of (9). For this we show that in each 2 species system the orbits of (9) converge to **H** (For the one species subsystems this is obvious.) If a two species subsystem of (11) has an internal equilibrium, say $\mathbf{E_{12}}$ then by assumption (ii),this is not a saddle but a sink, and hence $a_{11}a_{22} > a_{12}a_{21}$. This implies that the 1-2 submatrix of A is VL-stable [12,ch.15.3,],and [45] implies that $\mathbf{E_{12}}$ is the global attractor for the spatially structured system (9) restricted to the first two species.

If a two species subsystem of (11) has no internal equilibrium, then a one species equilibrium say $\mathbf{E_1}$ is the global attractor. Then Lemma 2 (see below) show the corresponding result for (9).A similar proof applies to the case where the origin **0** is the global attractor of a two species subsystem.

Finally observe that under the above assumptions (9)-(10), the invasion rates $\lambda_i$ at spatially homogenous boundary equilibria are the same for (9) and (11) Since all boundary equilibria are unsaturated for (11), they are for (9) as well. The same applies in case (iv) to the heteroclinic cycle. Note that the condition (12) on the eigenvalues is equivalent to the fact that all invariant measures supported on the heteroclinic cycle are unsaturated, see [16,Ex.4.5]

Species we mean types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

**Lemma 2:** Suppose, in a two species Lotka –Volterra system (11) (k=2) all interior orbits converge to the one species equilibrium $\mathbf{E_1}$ .Then all orbits of the spatial version (9) (k=2) converge to the spatially homogeneous one species equibrium $\mathbf{E_1} \in \mathbf{H}$.

**Proof:** If the system is of predator-prey or mutualistic type then Hasting's [45] result applies again. So let us assume that the local interaction is given by the competition system

(13) $\frac{dx_1}{dt} = r_1x_1(1 - x_1 - ax_2)$ $\frac{dx_2}{dt} = r_2x_2(1 - \beta x_1 - x_2)$

With 0<a <1< $beta$.In this case, Kazuo Kishimoto (letter to JH, Oct 1987) has given the following argument. There exists a family of forward invariant rectangles contracting to $\mathbf{E_1}$. This shows that all interior solutions of (9) (k=2) converge to $\mathbf{E_1}$.

Theorem 4 shows that a spatial network of identical patches is permanent, if the within patch dynamics is permanent, and three species Lotka-Volterra dynamics. It is not clear how to extend this result to more than three species. The crucial step in the three species case is that permanence precludes bistable two species subsystems. The spatial version of a bistable two species system allows plenty of stable equilibria outside **H,** see [33].However, permanent four species Lotka Volterra systems may have two species bistable subsystems [12,ch.16.4].So there may be invariant sets on the boundary outside **H** which need to be unsaturated to make (9) permanent.

## 6. MORE APPLICATIONS

6.1. **Disease, particle, product, marketing, financial, econometric dynamics with density-dependent demography**. GAO and Hethcote [46] introduced a model of disease dynamics with density-dependent demography. Here we consider a variation of their model in which a population of size P has a per-capita birth rate b(P) and per –capita death rate d. To ensure that the population persists in the absence of the disease, we assume that there exists K>0 such that b(K)=d. To describe the disease dynamics, let *S* be the number of individuals susceptible to the disease, *I* the Number of infected individuals, and R the number of individuals that have recovered from the disease. We assume that P=S+*I* +R. On the same lines we can think of sick and non performance accounts, and those which are conducting themselves well, and those which are satisfactory and find a parallel formula. The same thing holds true in case of a black hole eating matter and black hole radiating radiation.. If the disease transmission is asymptotic, β is the contact rate between susceptible and infected individuals, γ is the rate at which individuals recover from the disease, and m the mortality rate due to disease, then the population dynamics is given by

$$\frac{dp}{dt} = b(P)P - dP - mI$$

$$\frac{dI}{dt} = \beta(P - I - R)\frac{I}{\epsilon + P} - (\gamma + d + m)I$$

$$\frac{dR}{dt} = \gamma I - dR$$

Where ϵ>0 is a constant It is useful to introduce a change of coordinates in which $y=\frac{I}{P}$ and $z=\frac{R}{P}$:

$$\frac{dP}{dt} = (b(P)-d-my)P$$

$$\frac{dy}{dt} = \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(P)\ y$$

$$\frac{dz}{dt} = \gamma y - (b(P) - my)z$$

Setting $x_1$=P and $x_2$=(y,z) 'where' denotes transpose yields a structured model where $A_1(x)$=b-d-my and

$$A_2(x) = \begin{pmatrix} \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(p) & O \\ \gamma & my-b(p) \end{pmatrix}$$

Since {0,(*K*,0,0)}is a Morse decomposition of the boundary dynamics Theorem 3 and Remark 1 imply that this model is robustly permanent if and only if b(0)>d and $\beta\frac{K}{\epsilon+K} > \gamma + m + d$. In particular, for $\epsilon > 0$ sufficiently small, one requires that the basic reproductive number $\frac{\beta}{r+m+d}$ is greater then one.

Remarkably this same criterion determines robust permanence of model with significantly more complicated boundary dynamics. For example, suppose the focal population is a predator species. If the prey has abundance *N* and exhibits logistic dynamics f(*N*)is the per-capita predator consumption rate of the prey and b(N)is the per-capita reproductive rate of the predator then the dynamics become

$$\frac{dN}{dt} = rN(1 - N/K - f(N)P$$

$$\frac{dP}{dt} = (b(N) - d - my)P$$

$$\frac{dy}{dt} = \left( \beta(1 - y - z)\frac{P}{\epsilon+P} - r - m(1 - y) - b(N)\ y \right)$$

$$\frac{dz}{dt} = \gamma y - (b(N) - my)z$$

Where r>0 and *K*>0, and *f*(*N*),b(N)satisfy f(0)=b(0)Define $x_1 = N$, $x_2 = p$, $x_3 = (y,z)'$ Let us assume that the disease-free community (i.e.y=z=0)is permanent (i.e.b(K)>d) and let A be the global attractor in the interior of the P-N plane Under these assumption a Morse decomposition of the boundary dynamics is given by $(x_1, x_2, x_3) = (0,0,0)$, $(x_1, x_2, x_3) = (K, 0,0)$and A. Since the equilibria (0,0,0) and (k,0,0) are unsaturated, it remains to characterize $\lambda_3(\mu)$for any invariant measure μ supported by A. Proposition 3 and Remark 1 imply

$$\lambda_2(\mu) = \int_c (bx_1)\text{-}d)d\mu(x) = 0$$

*Hence* $= \int_c bx_1)d\mu(x) = d$ *and*

$$\lambda_3(\mu) = \beta \int \frac{x_2}{\epsilon+x_2} d\mu(x) - r - m - d$$

**Theorem 3**: and Remark 1 implies that this structured model is robustly permanent if $\frac{\beta}{r+m+d} > 1 and\ \epsilon > 0$ is sufficiently small.

6.2. Gene networks. The repressilator is an oscillatory gene network based on three (or more generally an odd number of transcriptional repressors. Such a system has been genetically engineered in E-coli by Elowitz and Leibler [47] Mathematical model go back to [48.49] they all involve concentrations of proteins and mRNs. A modified model with auto-activation was suggested in [50]

$$(14)\ x_i = \beta(y_i - x_i)$$

$$(15)\ y_i = aF(x_{)i} - y_i)$$

With

$$(16)\quad F(x)_i = x_i g\ (x_i, x_i - 1) = \frac{x_i}{1+x_i+p x_{i-1+Kp\, x_i x_i - 1}}$$

Here $x_i$ and $y_i$ are normalized concentrations of proteins and mRN As belonging to gene $i$,p is the strength of repression and k a cooperativity parameter this is a structured system of type(2)

$$(17)\ \begin{pmatrix} x_i \\ y_i\ dot \end{pmatrix} = \begin{pmatrix} -\beta & \beta \\ ag\,(x_i, x_i) & -1 \end{pmatrix} \begin{pmatrix} x_i \\ y_i \end{pmatrix}$$

In [50,Thm 3] it was shown that for n=3,a>1 and p>1,system(14)has a heteroclinic cycle connecting the 3 single gene equilibria (similar to the rock –scissors-paper dynamics of section5.2)Moreover

- the system is permanent i.e. the boundary of $R^6$+is repelling, if $\lambda + \mu > 0$
- this heteroclinic cycle is asymptotically stable if $\lambda + \mu < 0$

where $\lambda > 0$ and $\mu < 0$ are the invasion rates of the two missing genes i.e. the leading eignvalues of the 2X2 matrix in (17)evaluated at the appropriate single gene equilibria (for explicit expressions see[50,eqs.(127,128)].)

Theorem 3 implies the permanence condition implies the stronger conclusion of robust permanence A similar result holds for n odd, but the heteroclinic cycle is then actually a heteroclinic network.

Note: For references please see the original paper.

One of the examples for which we apply the model are as follows: **Progress in Particle and Nuclear Physics Volume 88, May 2016, Pages 1–28 Chiral magnetic and vortical effects in high-energy nuclear collisions—A status report D.E. Kharzeeva, b, J. Liaoc, d, , , S.A. Voloshine, G. Wangf**

The interplay of quantum anomalies with magnetic field and vorticity results in a variety of novel non-dissipative transport phenomena in systems with chiral fermions, including the quark–gluon plasma. Among them is the Chiral Magnetic Effect (CME)—the generation of electric current along an external magnetic field induced by chirality imbalance. Because the chirality imbalance is related to the global topology of gauge fields, the CME current is topologically protected and hence non-dissipative even in the presence of strong interactions. As a result, the CME and related quantum phenomena affect the hydrodynamical and transport behavior of strongly coupled quark–gluon plasma, and can be studied in relativistic heavy ion collisions where strong magnetic fields are created by the colliding ions. Evidence for the CME and related phenomena has been reported by the STAR Collaboration at Relativistic Heavy Ion Collider at BNL, and by the ALICE Collaboration at the Large Hadron Collider at CERN. The goal of the present review is to provide an elementary introduction into the physics of anomalous chiral effects, to describe the current status of experimental studies in heavy ion physics, and to outline the future work, both in experiment and theory, needed to eliminate the existing uncertainties in the interpretation of the data. Keywords Chiral magnetic effect; Chiral vortical effect; Chiral anomaly; Quark-gluon plasma; Heavy ion collisions Corresponding author at: Physics Department and Center for Exploration of Energy and Matter, Indiana University, 727 E Third Street, Bloomington, IN 47405, USA. Copyright © 2016 Elsevier B.V. All rights reserved.

## Sufficient Conditions For The Existence And Linear Stability Of Time Independent Age Distributions In Respect Of Particles, Propositions, Problems, Hegelian Structures, Deleuzean Treatises, Interactions And Transactions In Cosmos And Others

**Abstract:** Gopalaswamy studied sufficient conditions are obtained for the existence and linear stability of time independent age distributions in two species competition with age time lagged density dependent mortality and fertility functions.. Work was published in Bulletin Australian Mathematical Society Volume 25 (1982). We show how the model could be applied in case of finance, econometrics, psycho economics, interactions between particles and transactions in nature. Basically the competition is for resources, notwithstanding such experiments could be conducted in laboratory like a rat kept in front of a cat, such incidents could be represented mathematically, and the corresponding linear stability conditions ca be obtained in case of interactions in cosmos. Take for example the case of a cash credit account of a borrower. Many inward bills waiting for paid, or inland or foreign bills payment all are in line for the debit of the account, which also acts as a resource. In addition there are recurrent cases where in the credit is issued to the cash credit account in case of demand and inland bills. In case of economics, investment has inventories and it has to be utilized towards the end of consummation of output. And this is case of output gobbling part of investment which is inventories. Output has to be sold to the customer and the amount received is diverted partly to savings and consumption and part could be employed towards the consolidation of the system. So this is a case of investment itself acting as a both receiver and payer. It shall eat and also produce the finance. Savings as stated above could be deployed for the further diversification and proliferation and that is towards the end of putting the company on terra firma. Another example in cosmology is the case of black holes guzzling matter. There are lot

of black holes and lot of matter. Result of such a guzzling the production of Hawking radiation. Horava gravity is another example. How the conformal field theory could be changed in the case of Horava gravity. Still another case is the case of Vonneumann algebra in the cae of Horava gravity. In this case like cash credit account, the Horava gravity and Von Neumann algebra act as both receivers and takers. Such incidents are legion in cosmology and even in philosophy. We have considered self as individual general ledger, and the debits and credits are raised to the same account thus anointing as both a receiver and payer. We give plethora of such cases and apply the model. Due to spatial constants, same module is used for all the modules and we rewrite to state that the model is applicable to all modules. This is due to spatial restrictions. What is to be done is that each and every parameter is to be changed with a subscript or a superscript and finally it is to be concatenated towards the end of obtention of prediction model or optimizing model, for Duality of motion model for further analysis. Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would

have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

**Introduction**

Let $P_1$ *(t)* and $P_2$ *(t)* denote the total particles, interactions, transactions, etc, which we dilate upon further investigation. Perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. **Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems** sizes (or biomasses) at time $t \geq 0$ of two interacting species living in a common habitat and competing for a common pool of resources. The competition for resources will be implicit in our model similar to that in the two species Lotka- Volterra competition system. Assuming constant sex ratios in the two species we can consider $P_1$ *(t)* and $P_2$ *(t)*to be the Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

population of f; immigration, emigration and internal dispersion in the habitat are assumed to play no significant role in the dynamics of the community.Popoulation we mean Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

We suppose that the two species contain respectively $p_1(a, t)\,da$ and $p_2(a, t)\,da$ individuals with ages between $a$ and $a + da$ $(a \geq 0)$ at time $t$ so that we have

(1.19) $P_1(t) = \int_0^\infty \rho_1(a, t)\,da,\ P_2(t) = \int_0^\infty \rho_2(a, t)\,da,\ t \geq 0.$

The rates of change of the two species densities are by definition $D\rho_1(a, t)$ and $D\rho_2(a, t)$ where

(1.20) $D\rho_i(a, t) = \limsup_{h \to 0} \frac{\rho_i(a+h, t+h) - \rho_i(a,t)}{h},\ i = 1, 2.$

Assuming the existence of the limits in (1.2) we consider the following time lagged model system

(1.21) $$\begin{cases} D\rho_1(a,t) + f_1(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_1(a, t) = 0, & a>0,\ t>0 \\ D\rho_2(a,t) + f_2(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_2(a, t) = 0, & \text{'} \\ \rho_1(0, t) = \int_0^\infty b_1(a, P_1(t), P_2(t-\tau))\rho_2(a, t)da, & t>0, \\ \rho_2(0, t) = \int_0^\infty b_2(a, P_1(t), P_2(t-\tau))\rho_1(a, t)da, & \end{cases}$$ (1.4)

(1.5) $\rho_i(a, s) = \varphi_i(a, s),\ a.0,\ s \epsilon [-\tau, 0],\quad i = 1, 2,$

(1.6) $P_i(s) = \int \varphi_i(a, s)\,da,\ s \in [-\tau, 0],$

Where $\tau$some is fixed nonnegative constant; $f_1$ and $f_2$ denote the age and density dependent mortality rates with time lags in interspecific interactions, $b_1$ and $b_2$ denote the age and density dependent fertility functions again having time lags in the interspecific interactions; $\varphi_1$ and $\varphi_2$ denote the initial age distributions needed for the formulation of the model. Note we mean the population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

The model system (1.1)-(1.6) will be meaningful only if $\varphi_1, \varphi_2, f_1, f_2, b_1, b2$ are nonnegative functions of the respective arguments. Also since the total initial populations have o be finite, $\varphi_1(\cdot, s)$ should belong to $L_1(R^+)$ for each $s \epsilon$ $[-\tau, 0]$. That $f_i$ and $b_i$ $(i = 1, 2)$ are independent of t explicitly and this assumption can be interpreted to represent the temporal uniformity of the environment.

One of the fundamental questions for (1.1)-(1.6) is the following; under what conditions on $f_i, b_i$ and $\tau$ does the system (1.1)-(1.4) have time independent nontrivial solutions, and when such solutions, and when such solutions exist, are they stable with respect to some suitable stability criterion? When such stable time invariant solutions exist for (1.1)-(1.4) we say that the two types has a stable age distribution. Existence of a stable age distribution is the analogue of the existence of a stable steady state for the corresponding age independent system. Types we mean

types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

In this article we investigate the above question for the case of a competitive interaction with some additional constraints on the vital rates $f_i$ and $b_i$ $(i = 1, 2)$. Age dependent population systems without time delays have been considered by several authors (Gurt in and MacCamy [3, 4, 5], Gurt in and Levine [6], Haimovici [7] and Rotenberg [8]). The question of the relation between an age independent system and a corresponding age dependent system has been considered by Gurt in and MacCamy [4] and Gurt in and Levine [6] who have established an asymptotic (*as* $t \to \infty$ ) relation between such models by constructing a higher dimensional 1 umped parameter system (in terms of ordinary differential equations) to represent the age dependent distributed parameter system. Assuming the existence of stationary age distributions Haimovici [7] considers their stability in a system of two interacting populations explicitly taking into consideration the dynamical nature of the encironments resources and pollution.

In this article we add another realistic feature namely time lags in the age dependent models and obtain sufficient conditions for the existence of stationary age distributions in (1.1)-(1.4) and show that the time lags in interspecific interactions have no effect on the linear stability of the age distributions although the decay rates of perturbations will depend on the magnitudes of the time delays. Such "harmless" natures of time lags in interspecific interactions have been noted by the author for the age independent models in [1, 2].

## 7. Existence and uniqueness of solutions

We first narrate our assumptions on the vital rates of $f_i$ and $b_i$ $(i = 1, 2)$ so that the system (1.1)-(1.6) will denote a competition system. Let $BC\ (0, \infty)$ denote the linear space of bounded continuous real functions on $(0, \infty)$ with the norm defined by

$$||f|| = \{max\ |\ f(t)\ |,\ t \in (0, \infty)\}.$$

Let $K_1$ be the cone of nonnegative functions of $BC\ [0, \infty]$. Let $C\ [-\tau, 0]$ denote the space of real continuous functions with a norm defined by $||g|| = \{\ max\ |g(s)\ |,\ s \in [-\tau, 0]\}$; let $K_2$ be the nonnegative cone of functions in $C\ [-\tau, 0]$ . We present our first set of assumptions on $f_i$, $b_i$ and $\varphi_i$ for $i = 1, 2$.

$A_1$. $f_1 \in C\ (R^+ \times K_1 \times K_2, R^+)$, $f_1 \geq d_1^*$ on $R^+ \times K_1 \times K_2$

$f_2 \in C(R^+ \times K_2 \times K_1, R^+)$, $f_2 \geq d_2^*$ on $R^+ \times K_2 \times K_1$,

( $d_1^*$, $d_2^*$ are positive constants);

$f_1$ and $f_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ ; that is, *for* $i = 1, 2$,

$$|f_i(a, P_1', P_2') - f_i(a, P_1'', P_2'')| \leq d_i [||P_1' - P_1'', P_2' - P_2''||]$$

($d_1$, $d_2$ are positive constants),

$$\partial f_1/\partial P_1, \partial f_1/\partial P_2 \in C(R^+ \times K_1 \times K_2, R^+),$$

$$\partial f_2/\partial P_1, \partial f_2/\partial P_2 \in C(R^+ \times K_2 \times K_1, R^+),$$

$A_2$. $b_1 \in C(R^+ K_1 \times K_2, R^+)$, $b_1 \leq \beta_1^*$ *on* $R^+ \times K_1 \times K_2$,

$b_2 \in C(R^+ K_2 \times K_1, R^+)$, $b_2 \leq \beta_2^*$ *on* $R^+ \times K_2 \times K_1$,

*($\beta_1^*$, $\beta_2^*$ are positive constants);*

$b_1$, $b_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ *; that is for* $i = 1, 2$,

$|b_i(a, P_1', P_2') - b_i(a, P_1'', P_2'')| \leq \beta_i [||P_1' - P_1''|| + ||P_2' - P_2''||]$

$B_1$, $B_2$ Are positive constants),

$b_i(\circ, (P_1, P_2)) \epsilon L^{\infty}(R^+, R^+)$ For fixed $P_1, P_2, i = 1,2$,

$$\frac{\partial b_1}{\partial P_2}, \frac{\partial b_1}{\partial P_2} \epsilon\ C\ R^+ \text{x}\ k_1 \text{x}\ k_2, \text{R}^-),$$

$$\frac{\partial b_2}{\partial P_1}, \frac{\partial b_2}{\partial P_2} \epsilon\ C\ R^+ \text{x}\ k_2 \text{x}\ k_1, \text{R}^-),$$

$A_3$. $\varphi_i(\cdot, s) \in L^2(R^+, R^+) \cap C(R^+, R^+)$ for each $s \in [-\tau, 0]$,

$i, = 1, 2,$

By a solution of $(1.1) - (1.6)$ on $[0, T]$, we mean

$P_i: R^+ \times [-\tau, T] \to R^+$ $(i, = 1, 2,)$ With the following properties:

$S_1$. $P_i(\cdot, t) \in L^1(R^+, R^+)$ for each $t \in [-\tau, T]$;

$S_2$. $P_i(a, \cdot) \in C([-\tau, T], R^+)$ for each $a \in R^+$;

$S_3$. $DP_i$ Exists along the characteristics $t - a =$ constant

On $R^+ \times \{R^+ \cap [0,T]\}$ And is continuous $(i, = 1, 2,)$;

$S_4$. $P_i$ Satisfies the system (1.1)- (1.6) for

$(a,t,) \in R^+ \times [-\tau, T]$ .

Let us first convert the system (1.1)- (1.6) into an equivalent system of integral equations. For brevity we denote by P the pair ( $P_1, P_2$) If we let

$$P_i(a+x, t+x) = P_i(x), (i, = 1, 2,),$$

Then (1.3) considered at $(1.3)$ considered at $(a+x, t+x)$ becomes

$$\frac{dp_i(x)}{dx} + f_i(a+x, P(t+x)) \; P_i(x) = 0$$

Which has a solution of the form?

$$P_i(x) = P_i(0) \exp\left[ - \int_0^\infty f_i(a+s, P(t+S)) ds \right]$$

And hence

(2.1) $p_1(a+x, t+x) = p_1(a,t) \, exp\left[ - \int_o^x f_1(a+s, p(t+s)) ds \right]$

For $a \geq t$ *we choose* $(a,t)=(a-t,0)$ in (1.3) and $x=t$ in (2.1) so that

(2.2) $p_1(a,t) = \varphi_1(a-t, 0) \exp\left[ - \int_o^t f_1(a-t+s, p_1(s)) ds \right]$

$$O \leq t \leq a, i = 1,2,$$

With

$$P_1(s)=(p_1(s),p_2(s-t)),$$

$$P_2(s)=(p_1(s-t),p_2(s)),$$

For $0 \leq a < t$ we choose $(a,t)=(0,t-a)$*in (1.3)and* $x=a$ *in (2.1)* so that

(2.3) $p_i(a,t) = B_i(t-a) exp[-\int_o^a f_1(s, p_i(t-a+s)) ds]$

$$t > a \geq, i = 1,2,$$

Where

| |
|---|
| (2.4) $B_i(t) = p_i(0,t), i = 1,2, t > 0.$ |
| Define $M_i$ and $L_i$ $(i = 1,2)$ as follows: |
| $M_i(a,t,p_i(t)) = exp[-\int_o^{t-a} f_1(s, p_i(a-a+s))ds]$ |
| (2.5) $L_i(a,t,p_i(t)) = \exp[-\int_o^t f_1(a+s, p_i(s))ds]$ |
| It will now follow from (1.1),(1.3),(2.2)-2.5)that |
| (2.6) $p_i(t) = \int_o^t B_1(a) M_1(a,t,p_i(t)da + \int_o^\infty \varphi_i(a,0) L_i(a,t,p_i(t)da,$ |
| $t > 0, i = 1,2,$ |
| (2.7) $p_i(s) = \Phi_i(s) = \int_o^\infty \varphi_i(a,s)da, s \in [-\tau, 0], i = 1,2,$ |
| (2.8) $\beta_i(t) = \int_0^t \beta_i(a) b_i(t-a, P_i(t)) M_i.(a,t,P_i(t))da$ |
| $+\int_0^\infty b_i(a+t, P_i(t))\varphi_i(a,0)L_i(a,t,P_i(t))da$ |
| $i = 1, 2, t > 0$ |
| the equivalence of (2.6)-(2.8) with (1.1)-(1.6) is established by the following whose proof is identical to a similar result of Gurtin and MacCamy [4]; hence we omit the details of the proof of the following. |
| THEOREM 1 : *Suppose the system* (1.1)-(1.6) *satisfies the conditions* $A_1, A_2, A_3$ . *If* $(P_1, P_2)$ *is a solution of* |
| *(1.1)-(1.6) then the total population* $P_1, P_2$ *and the birth rates* $B_1, B_2$ *satisfy the integral equation (2.6)-(2.8).* |
| *conversely if* $P_1, P_2, B_1, B_2$ *are nonnegative continuous solution of (2.6)-(2.8) on [ 0, T] and if* $P_1, P_2$ *are* |
| *defined by (2.2)-(2.4) on* $R^+$ x [0, T] *then such* $P_1, P_2$ *provide a solution of (1.1)-(1.6) on* $R^+$ x [0, T] . |
| The following a *priori* estimates are useful to prove our existence theorem below. It will immediately follow from the bounds for $f_i$ and $b_i$ that |
| |
| $M_i(a,t,P) \leq \exp[-d_i^*(t-a)] \leq 1, if \quad t \geq a,$ |
| *(2.9)* |
| $L_i(a,t,P) \leq \exp[-d_i^* t]$ |
| From (2.7)-(2.8) we derive that |

$$\beta_i(t) \le \int_o^t \beta_i^* \beta_i(a) \exp[-d_i^*(t-a)]\, da + \beta_i^* \Phi_i(0) \exp[-d_i^* t], \quad i = 1,2,$$

And hence by Gronwall's inequality,

(2.10) $$\beta_i(t) \le \beta_i^* \Phi_i(0) \exp[\delta_i t], \ \delta_i = \beta_i^* - d_i^*$$

(2.5), (2.6) and (2.10) lead to

(2.11) $$|P_i(t) - 0_i(0)| \le \beta(0)\, exp[\delta_i\, a - d_i^*(t-a)]\, da$$

$$+\int_0^\infty \varphi_i\ (a,0)|1 - exp(-d_i^* t)|da$$

$$\le \beta_i^* \Phi_i\ (0)\, exp(\beta_i^* t) \left\{\frac{1-exp[-(d_i^*+\beta_i^*)t]}{d_i^*+\beta_i^*}\right\} + \beta(0) d_i^* t$$

$$\le [\beta_i^* \beta(0)\, exp(\beta_i^* t) \beta(0) d_i^*]t \quad for\ t > 0.$$

We can now prove the following result on the existence and uniqueness of solution of the integral equations (2.)-(2.8)

**THEOREM 2:** *Assume that the system* (1.1)-(1.6) *satisfies the hypotheses*

$A_1, A_2, A_3$. *Then there exists a unique set of continuous function* $(P_1, P_2, B_1, B_2)$ such *that*

$$P_i : R^+ \to R^+, \qquad B_i : R^+ \to R^+ \ (i = 1,2),$$

*Which satisfy the integral equations* (2.6)-(2.8)

Proof. We will first show the existence of local solution on same (possibly small) interval [0, T] $\subset R^+$ such that

$$P_i \in C([0,T], R^+), \quad B_i \in C([0,T], R^+), i = 1,2,$$

Satisfying (2.6)-(2.8) and then show the $P_i$ and $B_i$ can be continued as solution for any finite interval in $R^+$ for some fixed

Positive number $T$ we let

$$C^+[-\tau, T] = \{f \in C([-\tau, T], R) | f \ge 0\}$$

Let $||.||_T$ denote the supremum norm in $C^+[-\tau, T]$. Then for any

$(x, y) \in C^+[-\tau, T] \times C^+[-\tau, T]$ we define

$||(x,y)|| = ||x||_T + ||y||_T,$

$C_T = C^+[-\tau, T], \quad C_T^2 = C^+[-\tau, T] \times C^+[-\tau, T].$

Members of $C_T$ will be denoted by $Z_t$ with the meaning that

$Z(t+s) = z_t(s)$ *for* $s \in [-\tau, T]$ *and* $t \in [0, T]$. Now for any fixed

$P_t = (P_{1t}, P_{2t}) \in C_T^2$, (2.8) is a linear system of uncoupled Volterra integral equations in B = $(B_1, B_2)$ and hence (2.8) will

have a unique continuous solution under the assumption $A_1 - A_3$. let us denote such a solution by

$$B(t) = B^T(P_t), \quad t \in [0, T]$$

Or

$$B_1(t) = B_1^T(P_{1t}, P_{2t}), \quad B_2(t) = B_1^T(P_{1t}, P_{2t}), \quad t \in [0, T]$$

If we supply this solution $(B_1, B_2)$ in (2.6) we find that (2.) will be satisfied if and only if $(P_{1t}, P_{2t})$ is a fixed point of the

operator

$$\pi(P_t) = \{\pi_1(P_{1t}, P_{2t}), \pi_2(P_{1t}, P_{2t})\}$$

Where $\pi_1$ *and* $\pi_2$ respectively denote the right sides of (2.6) for *i =1,2* .

It is not difficult to see from the nature of $M_i$, $L_i$ and the hypotheses on $f_i, b_i, \varphi_i$ that

$$B: C_T^2 \to C_T^2 \text{ and } \pi: C_T^2 \to C_T^2$$

(which guarantees the nonnegativity of $P_i$ and $B_i$). We will first show that

There exist positive constants $r^*, t^*$ such that $\pi$ is a mapping of a sphere

$$S_{t*}(\Phi, r^*) = \{P_t = (P_{1t}, P_{2t}) \in C_{t*}^2 \mid \|P_t - \Phi_t\|_{t*} \le r^*\}$$

Into itself and also is a contraction where $\Phi_t$ is defined on [$-\tau$, $t^*$]

By the following:

$$\Phi_t = (\Phi_{1t}, \Phi_{2t});$$

(2.12)

$$\Phi_{it} = \begin{cases} \int_0^\infty \varphi_i(a,s)\,da \ for\ s \in [-\tau, 0]\,. \\ \int_0^\infty \varphi_i(a,0)\,da \ for\ s \in [0, t^*]\,. \end{cases}$$

We note that similar to (2.11) one gets the estimates

$$|\mathcal{P}(P)(t) - \Phi_i(0)| \le \{B_i^*\Phi_i(0)\exp[B_i^* t] + \Phi_i(0)d_i^*\}t$$

Which together with (2.7) and (2.12) lead to

(2.13) $$\|\mathcal{P}(P_t) - \Phi_t\|_{t*} \le \begin{cases} \sum_{j=1}^2 \|\Phi_{jt}\|_{t*}[B_j^*\exp(B_j^* t) + d_i^*] \\ t^* \end{cases}$$

Thus if we choose,

$$r^* = T\sum_{j=1}^2 \|\Phi_{jt}\|_{t*}[B_j^*\exp(B_j^* T) + d_i^*]$$

Then $\mathcal{P}$ is a mapping of $S_{t*}(\Phi, r^*)$ into itself for all $t^* \le T$. we then will have to show that for small $t^* > 0$, $\mathcal{P}$ is a contraction.

Let $P_t^{(1)}$ *and* $P_t^{(2)}$ be arbitrarily chosen in $S_{t*}(\Phi, r^*)$ and consider $|\mathcal{P}(P_t^{(1)}) - \mathcal{P}(P_t^{(2)})|$ for t $\le t^*$.

Corresponding to the chosen $P_t^{(1)}$, $P_t^{(2)}$ we let $B^{(1)}(a) = B(P_T^{(1)})$ and $B^{(2)}(a) = B(P_T^{(2)})$. Then from

(2.6),

$$|\mathcal{P}(P^{(1)}(t)) - \mathcal{P}(P^{(2)}(t)| \le \left|\int_0^t B_1^{(1)}(a)\left\{M_1\left(a,t,P_t^{(1)}\right) - M_1\left(a,t,P_t^{(2)}\right)\right\}da\right|$$

$$+ \left|\int_0^t [B_1^{(1)}(a) - B_2^{(2)}(a)]\, M_1(a,t,P_t^{(2)})da\right|$$

$$+ \left|\int_0^t \varphi_1(a,0)\left[L_1\left(a,t,P_t^{(1)}\right) - L_1\left(a,t,P_t^{(2)}\right)\right]da\right|$$

$$= J_1 + J_2 + J_3 \text{ (say)}$$

using the elementary inequality

$$|e^x - e^y| \le |x - y| \ for\ x, y \le 0$$

We derive from (2.5) that

$$\left| M_1(a,t,\ p_t^{(1)}) - M_1(a,t,\ p_t^{(2)}) \right|$$

$$\leq \int_o^{t-a} \left| f_1(s,\ p_t^{(1)}(a+s)) - f_1(s,\ p_t^{(2)}(a+s)) \right| ds$$

$$\leq d_1 \| p_t^{(1)} - p_t^{(2)} \|_{t^*} \text{(using } A_1\text{)}$$

*Now,using (2.10),*

$$\left| J_1 \right| \leq \int_o^a \beta_1^* \Phi_1(0)\, e^{\delta 1^a} \left( \int_o^{t-a} d_1 \| P_t^{(1)} - P_t^{(2)} \| ds \right) da$$

$$\leq d_1\, \beta_1^*\, \Phi_1(0)\, t^2 \| P_t^{(1)} - P_t^{(2)} \| e^{\delta 1^t/2}$$

In a similar way

$$\left| J_3 \right| \leq \int_o^{\infty} \varphi(a,o) \left[ \int_o^t \left| f_1(a+s, P_t^{(1)}(s) - f_1(a+s, P_t^{(2)}(s)) \right| ds \right] da$$

$$\leq d_1 \Phi_1(0) \| P_t^{(1)} - P_t^{(2)} \| t$$

Now we have to estimate $\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right|$ in order to estimate $J_2$, from (2.8),

$$B_1^{(1)}(a) - B_1^{(2)}(a)$$

$$= \int_o^a \left[ B_1^{(1)}(\sigma) - B_1^{(2)}(\sigma) b_1 \left( a-\sigma, P_t^{(1)}(a) \right) M_1(\sigma, a, P_t^{(1)}(a) d\sigma + \int_o^a \left[ B_1^{(2)}(\sigma) \right. \right.$$

$$\text{X} \left[ b_1(a-\sigma, P_t^{(1)}(a) M_1(\sigma, a, P_t^{(1)}(a) - b_1 \left( a-\sigma, P_t^{(1)}(a) M_1 \left( \sigma, a, P_t^{(2)}(a) \right) \right] d\sigma \right.$$

$$+ \int_o^{\infty} \varphi(\sigma, o) \left[ b_1 \left( \sigma + a, P_t^{(1)}(a) \right) - b_1 \left( \sigma + a P_t^{(2)}(a) \right) \right], L_1 \left( \sigma, a, P_t^{(1)}(a) \right) d\sigma$$

$$+ \int_o^{\infty} \varphi(\sigma, o) b_1 \left( \sigma + a, P_t^{(2)}(a) \right) \left[ L_1 \left( \sigma, a, P_t^{(1)}(a) \right) - L_1 \left( \sigma, a, P_t^{(2)}(a) \right) \right] d\sigma$$

$$= R_1 + R_2 + R_3 + R_4 \text{ (say)}.$$

By the uniform Lipschitz continuity of $b_1$ in $A_2$ we get

$$\left| b_1 \left( a+\sigma, P_t^{(1)}(a) \right) - b_1 \left( \sigma + a, P_t^{(2)}(a) \right) \right| \leq B_1 \| P_t^{(1)} - P_t^{(2)} \| \text{ for } a \in [o, t^*]$$

And hence we have

$$|R_3(a)+R_4(a)|\leq\beta_1\Phi_1(0)||P_t^{(1)}-P_t^{(2)}||_t^{*}+\Phi_1(0)\beta_1 d_1||P_t^{(1)}-P_t^{(2)}||_t^{*a},$$

$$|R_2(a)|=[\int_0^a[B_1^{(2)}(\sigma)\{b_1(a-\sigma, P_t^{(1)}(a))-b_1(a-\sigma, P_t^{(2)}(a))\}M_1(\sigma, a, P_t^{(1)}(a))$$

$$+b_1(a,-\sigma, P_t^{(2)}(a))\{M_1(\sigma, a, P_t^{(1)}(a))-M_1(a-\sigma, P_t^{(2)}(a))\}]d\sigma]$$

$$+b_1(a-0, P_t^{(2)}(a))\{M_1(\sigma, a, P_t^{(1)}(a))-M_1(a-\sigma, P_t^{(2)}(a))\}]d\sigma]$$

$$\leq \int_0^a [\beta_1^{*}\Phi_1(0)e^{\delta 1^a}\beta_1\| P_t^{(1)}-P_t^{(2)}\|_{t^*}+\beta_1 d_1\| P_t^{(1)}-P_t^{(2)}\|_{t^*}(a-o)]do$$

$$\leq \beta_1^{*}\Phi_1(0)e^{\delta 1^a}\| P_t^{(1)}-P_t^{(2)}\|_{t^*}\{\beta_1 a+d_1 a^2/2\}$$

Thus

$$\left|B_1^{(1)}(a)-B_1^{(2)}(a)\right|$$

$$\leq \int_o^a \left| B_1^{(1)}(\sigma)-B_1^{(2)}(\sigma)\right| B_1^{*d}\sigma$$

$$+\Phi_1(0)\|P_t^{(1)}-P_t^{(2)}\|_{t^*}\left\{B_1+B_2 d_1 a+B_1^{*}\left(B_2 a+d_1 a^2/2\right)e^{\delta 1^a}\right\}$$

$$\leq \beta_1^{*}\int_o^a \left| B_1^{(1)}(\sigma)-B_1^{(2)}(\sigma)\right| do+\left(r^{*}+\|\Phi_t\|_{t*}\right)\| P_t^{(1)}-P_t^{(2)}\|_{t^*} g_1^{(a)}$$

Where $g_1(a)\rightarrow \beta_1$ as $a\rightarrow 0$. By Gronwall's inequality it will follow that

$$\left|B_1^{(1)}(a)-B_1^{(2)}(a)\right|\leq \left(r^{*}+\|\Phi_t\|_{t*}\right)\| P_t^{(1)}-P_t^{(2)}\|_{t^*}\{g_1(a)+0(a)\}$$

Using this estimate for $R_1(a)$ we get

$$|R_2(a)|\leq \left| \int_0^a[B_1^{(2)}(\sigma) B_1^{(2)}(\sigma),] b_1(a))-b_1(a-\sigma, P_t^{(1)}(a)) M_1(\sigma, a, P_t^{(1)}(a))do\right|$$

$$\left|B_1^{*}\left(r^{*}+\|\Phi_t\|_{t*}\right)\| P_t^{(1)}-P_t^{(2)}\|_{t^*}0(a) as\ a\rightarrow 0\right.$$

*Thus*

$$\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right| \leq (r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ \; P_t^{(1)} - P_t^{(2)} /\!/_t o(1) as \quad a \to o$$

And hence

$$\left| \pi_1(P_t^{(1)} - \pi_1\left(P_t^{(2)}\right) \right| \leq ((r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ \; P_t^{(1)} - P_t^{(2)} /\!/_t o(t) as \quad t \to o$$

Which together with a similar bound for $\left| \pi_2(P_t^{(1)} - \pi_2\left(P_t^{(2)}\right) \right|$ leads to

$$\| \; \pi \; (P_t^{(1)} - \pi \left(P_t^{(2)}\right) /\!/ \; \leq (r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ \; P_t^{(1)} - P_t^{(2)} /\!/_{t*} o(t^*)$$

As $t^* \to o$ showing that if $t^*$ is small enough, the mapping $\pi$ on

$s_{t^*}(\Phi, r^*)$ is a contraction. Hence the system (2.6)-(2,8) has a unique solution for $o \leq t \leq t^* for \; t^*$ sufficiently small and positive

Now if [0,T] is any finite interval on $R^+$, we can consider [0,T] as a finite union of intervals of length less than or equal to $t^*$ Since the estimate (2.11)is valid for any finite interval we can extend the solution of (2,6)-(2.8) form [0, $t^*$] to [0,T] and this completes the proof,

3. Stationary age distributions and their linear stability we will now establish the existence of time independent solution

$P_t^*, P_t^*$ Of (1.1)-(1.6): such solutions satisfy

$$(3.1) \; \frac{dp\, P_2^*(a)}{da} + f_i(a, P_1^*, P_2^*) \; P_i^*(a) = 0, \;\; i = 1, 2, a > 0,$$

$$(3.2) \; P_i^* = \int_o^\infty \; P_i^*(a) da, i = 1, \; 2,$$

$$(3,3) \; P_i^* = \int_o^\infty \; b_i(a, \; P_1^*, \; P_2^*) P_i^*(a) da, \; i = 1. \;\; 2$$

All solutions of (3.1)are of the from

$$(3.4) \; P_i^*(a) = P_i^*(o) \exp\left[ - \int_o^\infty \; f_i(s, P_1^*, P_2^*) ds \right], i = 1, \quad 2, a \geq 0,$$

Which together with (3.3) lead to

$$1 = \int_o^\infty \; b_1(a, P_1^*, P_2^*) \exp\left[ - \int_o^a \; f_1(s, P_1^*, P_2^*) ds \right] da,$$

$$(3.5) \; 1 = \int_o^\infty \; b_2(a, P_1^*, P_2^*) \exp\left[ - \int_o^a \; f_2(s, P_1^*, P_2^*) ds \right] da$$

Thus the existence of stationary solutions of (1.1)-(1.6) reduces to the existence of a pair of positive constants $P_1^*, P_2^*$ satisfying (3.5). We can now prove the following,

THEOREM 3. Assume that $f_1$ and $b_i$ satisfy ( $A_1$ ) and ( $A_2$ )and let

(3.6) $F_{i.}\ (P_1, P_2) = \int_o^\infty b_i(a, P_1, P_2) \exp\left[-\int_o^\infty f_i\,(s, P_1, P_2)ds\right] da,$

$i$ =1, 2, $P_1 \geq 0$, $P_2 \geq 0$,

**And suppose the following hold:**

(3.7) $A_4$. $F_{i.}\ (0,\ 0) > 1,\ i = 1, 2$:

$A_5$. *there exists a positive constant c for which*

$F_1$(C,0)< $F_{2.}$(C,0),

(3.8) $F_1$(0.C)< $F_{2.}$(0,C),

Then there exists a unique pair $(P_1^*, P_2^*)$ of real numbers $P_1^* > o$,

$P_2^* > 0$ such that

(3,9) $F_1(P_1^*, P_2^*) = 1 = F_2(P_1^*, P_2^*)$

And the unique nonegative solution of (3.1)-(3.3) is given by

(3.10) $P_2^*(a) = [P_2^* \exp[-f_i\,(s, P_1^*, P_2^*)ds])/(\int_o^\infty \exp[-\int_o^a f_i\,(s, P_1^*, P_2^*)ds]da),$

$i$ =1, 2, $a \geq 0$,

Proof Consider the elements of the surfaces defined by

Z= $F_1(P_1, P_2), z = F_2(P_1, P_2),\ P_1 \geq 0, P_2 \geq 0,$

In the ($P_1$,$P_2$,z) space. By hypothesis ($A_1$) we have $\partial F_i/\partial P_j < 0\ and$ , by ($A_2$), $\partial b_i/\partial P_j \leq 0$ for all

$P_1 \geq 0, P_2 \geq 0$, $i$ =1, 2,

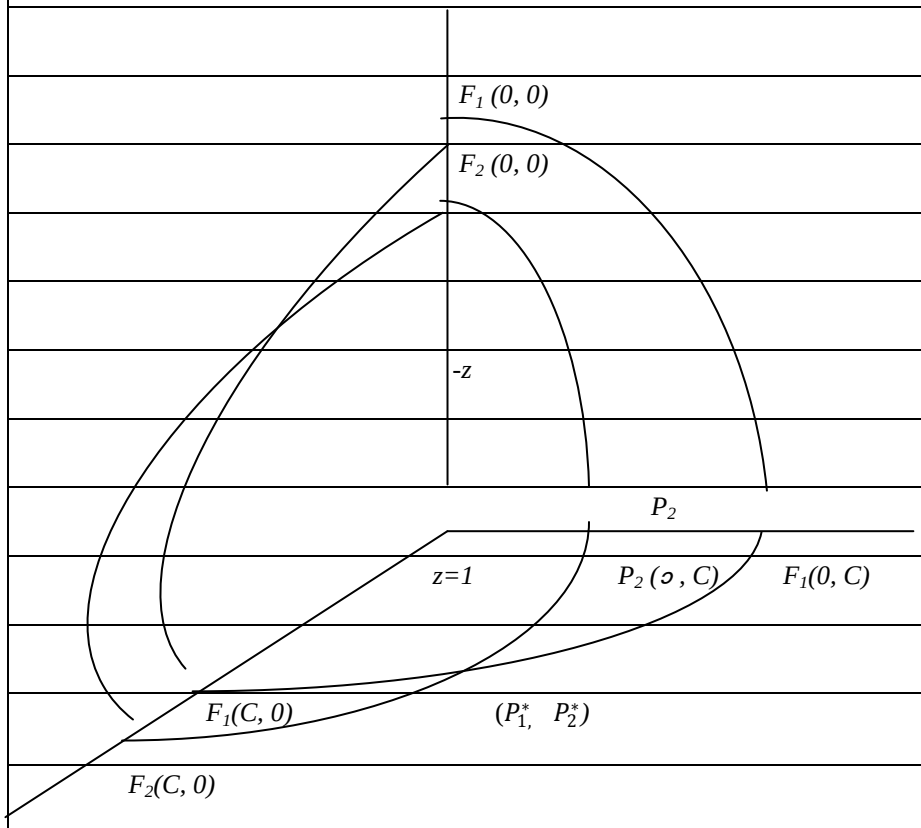
$j$ =1, 2, Hence by ($A_4$) the intersections of these surfaces with the plane z=1 define two curves on the z=1 plane connecting the I ines {z=1, $P_1 \geq 0, P_2 \geq 0,$} and {z= $1, P_2 \geq 0, P_1 \geq 0$, These two curves are defined by

$F_1(P_1, P_2 = 1\ and\ F_2(P_1, P_2) = 1,\ F_1(P_1, P_2 \geq 0,$

Which by ($A_s$) will intersect at a unique point say $((P_1^*, P_2^*)$ with $(P_1^* > 0, P_2^* > 0$ (see figure). The uniquess of the point $((P_1^*, P_2^*)$ *is a* consequence of the smoothness of the surfaces $z = F_1(P_1, P_2)$ and $z = F_2(P_1, P_2)$

Following is a hypothetical case for Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates attributable or ascribable in the transactions and reactionary potentialities.




(The figure corresponds to the case $F_1(0,0) > F_2(0,0) > 0$. The other possibility $1 < F_1(0,0)$ is treated similarly) It will now

Follow that for such a pair $(p^*_1, p^*_2)$ the unique nonnegative solution of (3.1) - (3.3) is given by (3.10).

Now to examine the linear stability of the stationary age distributions $p^*_1, p^*_2$, we let

(3.11)
$$\begin{cases} p_i(a,t) = p_i^*(a) + u_i(a,t) \\ p_i(t) = p_i^* + p_i(t), \quad i = 1, 2, \quad a > 0, \quad t > 0, \\ p_i(t) = \int_0^\infty u_i(a,t)da, \end{cases}$$

In (1.1)-(1.5) and derive the following variational system after neglecting the nonlinear terms in the perturbations $\boldsymbol{u}_i$ and $\boldsymbol{p}_i$ $(i = 1, 2,)$ (such a procedure of linear stability analysis can be justified As has been done in Gurtin and MacCamy [4].)

(3.12)
$$\begin{cases} Du_i(a,t) = -f_i\left(a,p^*_1,p^*_2\right)u_i(a,t) - p^*_i(a)\left\{\sum_{i=1}^{2}\frac{\partial f_i}{\partial p^*_{\square}}\,p_{i\square}(t)\right\}, \\ u_i(0,t) = \int_0^\infty b_i\left(a,p^*_1,p^*_2\right)u_i(a,t)da + \int_0^\infty p^*_i(a)\left\{\sum_{\square=1}^{2}\frac{\partial b_i}{\partial p^*_{\square}}\,p_{i\square}(t)\right\}da \end{cases}$$

Where $i = 1,2$ and (3.13)
$$p_{11}(t) = p_1(t),\quad p_{12}(t) = p_2(t-\tau),$$
$$p_{21}(t) = p_1(t-\tau),\quad p_{22}(t) = p_2(t)$$

To consider the asymptotic bahaviour as $t \to \infty$ as of solutions of (3.12) we let

(3.14)
$$\begin{cases} u_i(a,t) = u\square(a)\exp[\lambda t], \quad \square(a) \to 0 \text{ as } a \to \infty, \\ p_i(t) = p^*_i \exp[\lambda t], \\ \qquad\qquad i = 1,\ 2,\ a > 0,\ t > 0, \\ p^*_i = \int_0^\infty \square(a)da, \end{cases}$$

Using (3.14) in (3.12) - (3.13),

$$\frac{d\square}{da} + \lambda\square = -f_i\left(a,p^*_1,p^*_2\right)\square(a) - p^*_i(a)\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p^*_j}\,p^*_{i\square}\right\},$$

(3.15)

$$\square(0) = \int_0^\infty b_i\left(a,p^*_1,p^*_2\right)\square(a)da + \int_0^\infty p^*_i(a)\left\{\sum_{j=1}^{2}\frac{\partial b_i}{\partial p^*_j}\,p^*_{i\square}\right\}da,$$

Where $i = 1,\ 2,$ and

$$p^*_{11} = p^*_1,\quad p^*_{12} = p^*_2\exp\{-\lambda\tau\},$$

$$p^*_{21} = p^*_1,\ \exp\{-\lambda\tau\},\ p^*_{22} = p^*_2.$$

If we integrate the first of (3.15) and use the fact that $\square(a) \to 0$ as $a \to \infty$ we can derive that $p^*_1$ and $p^*_2$ are governed by

(3.16) $$\lambda p_1^* - \int_0^\infty p_1^*(a)\left[\left(\frac{\partial b_1}{\partial p_j^*} - \frac{\partial f_1}{\partial p_1^*}\right)p_1^* + \left(\frac{\partial b_1}{\partial p_2^*} - \frac{\partial f_1}{\partial p_2^*}\right)p_2^* e^{-\lambda\tau}\right]da$$

$$= \int_0^\infty \left[b_1(a, p_1^*, p_2^*) - f_1(a, p_1^*, p_2^*)\right]\varphi_i(a)da,$$

(3.17) $$\lambda p_2^* - \int_0^\infty p_2(a)\left[\left(\frac{\partial b_2}{\partial p_1^*} - \frac{\partial f_2}{\partial p_1^*}\right)p_1^* e^{-\lambda\tau} + \left(\frac{\partial b_2}{\partial p_2^*} - \frac{\partial f_2}{\partial p_2^*}\right)p_2^*\right]da$$

$$= \int_0^\infty \left[b_2(a, p_1^*, p_2^*) - f_2(a, p_1^*, p_2^*)\right]\varphi_2(a)da,$$

From the first of (3.15) we get

(3.18) $$\varphi_i(a) = \{\varphi_i(0)p_i^*(a)/p_i^*(0)\exp[\lambda a]\}$$
$$- p_i^*(a)\exp[-\lambda a]\int_0^a \left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}p_i^*\varphi_i\right\}e^{\lambda s}ds, \quad i = 1, 2,$$

Which by the second of (3.15) becomes,

(3.19) $$\varphi_i(a)\left[p_i^*(a)/p_i^*(0)c_i\exp(\lambda a)\right]\left[\int_0^a p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}p_i^*\varphi_i\right\}da\right.$$

$$\left. - \int_0^\infty b_i(a, p_1^*, p_2^*)p_i^*(a)e^{-\lambda a}\left\{\int_0^a\left(\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}p_i^*\varphi_i\right)e^{\lambda s}ds\right\}da\right]$$

$$- p_i^*(a)e^{-\lambda a}\int_0^a\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}p_i^*\varphi_i\right\}e^{\lambda s}ds, i = 1, 2,$$

where

(3.20) $$c_i = 11 - \frac{1}{p_i^*(0)}\int_0^\infty b_i(a, p_1^*, p_2^*)p_i^*(a)e^{-\lambda a}da, \quad i = 1, 2,$$

Now supplying $\varphi_1(a)$ and $\varphi_2(a)$ from (3.19) in (3.16)-(3.17) we derive that

$$(\lambda - A_{11})p_1^* - A_{12}p_2^*\exp(-\lambda\tau) = g_{11}(\lambda)p_1^* + g_{12}(\lambda)p_2^*\exp(-\lambda\tau),$$

(3.21)

$$-A_{21}p_1^*\exp(-\lambda\tau) + (\lambda - A_{22})p_2^* = g_{21}(\lambda)p_1^*\exp(-\lambda\tau) + g_{22}(\lambda)p_2^*$$

Where

(3.22) $$A_{ij} = \int_0^\infty p_i^*(a)\left[\frac{\partial b_i}{\partial p_j^*} - \frac{\partial f_i}{\partial p_j^*}\right]da,$$

(3.23) $g_{i}(\lambda) = [p_i^*(0)c_i]^{-1} \{\int_0^\infty \sigma_i(a)p_i^*(a)\, e^{-\lambda a}\, da\}$
$\left\{\int_0^\infty \left[p_i^*(a)\frac{\partial b_i}{\partial p_j^*} - b_i(a, p_1^*, p_2^*)p_i^*(a)\, e^{-\lambda a}\left(\int_0^a \frac{\partial f_i}{\partial p_j^*}\, e^{\lambda s}\, ds\right)\right] da\right\}$

$$-\int_0^\infty \sigma_i(a)p_i^*(a)\, e^{-\lambda a}\left(\int_0^a \frac{\partial f_i}{\partial p_j^*}\, e^{\lambda s}\, ds\right) da,\ i, j = 1, 2,$$

(3.24) $\sigma_i(a) = b_i(a, p_1^*, p_2^*) - f_i(a, p_1^*, p_2^*),\ i, = 1, 2.$

To solve the linear perturbational system (3.12) it is enough to find the constants $\lambda, p_1^*, p_2^*$ satisfying (3.21) since one can then use (3.19) to find the perturbations from (3.14). A nontrivial solution $(p_1^*, p_2^*)$ for (3.21) will exist if and only if $\lambda$ is a root of the equation $et\begin{bmatrix}\lambda - (A_{11} + g_{11}(\lambda)) & -(A_{12} + g_{12})e^{-\lambda\tau} \\ -(A_{21} + g_{21})e^{-\lambda\tau} & \lambda - (A_{22} + g_{22}(\lambda))\end{bmatrix} = 0$

Or equivalently

(3.25) $\lambda^2 - \lambda(A_{11} + A_{22}) + A_{11}A_{22} - A_{12}A_{21}\exp[-2\lambda\tau] = \lambda s_1(\lambda) + s_2(\lambda, \tau)$

Where $S_1(\lambda) = g_{11}(\lambda) + g_{22}(\lambda)$ ,

$S_2(\lambda, \tau) = [A_{21}g_{12}(\lambda) + A_{12}g_{21}(\lambda) + g_{12}(\lambda)g_{21}(\lambda)]\exp[-2\lambda\tau]$

$-A_{11}g_{22}(\lambda) - A_{22}g_{11}(\lambda) - g_{11}(\lambda)g_{22}(\lambda)$

The following result establishes the linear local asymptotic Stability of the stationary age distributions $p_1^*(a)$ and $p_2^*(a)$

THEOREM 4 Assume the following:

$A_6$. $0 < \beta_i^* < d_i^*\ (i, = 1, 2.)\ (see(A_1)\ and\ (A_2))$;

$A_7$. $\{A_{11} + A_{22} + \| S_1(0) \|\}^2 < 4\{(A_{11}A_{22} - A_{12}A_{21}) - \| S_2(0) \|\}$,

$\| S_1(0) \| = |g_{11}(0)| + |g_{22}(0)|$ ,

$\| S_2(0) \| = |A_{11}g_{12}(0)| + |A_{12}g_{21}(0)| + |g_{12}(0)g_{21}(0)|$

$+ |A_{11}g_{22}(0)| + |A_{22}g_{11}(0)| + |g_{11}(0)g_{22}(0)|$

Then all the roots of (3.25) have negative real parts and hence the

perturbations $u_i(a, t) = \pi_i(a)\exp[\lambda t] \to 0\ as\ t \to \infty$ for $i = 1, 2\ and\ a \geq 0$

Proof . Define $F_1(\lambda,\tau)$ and $F_2(\lambda,\tau)$ as follows:

$F_1(\lambda,\tau) = \lambda^2 - \lambda(A_{11}+A_{22}) + A_{11}A_{22} - A_{12}A_{21}\exp[-2\lambda\tau];$

$F_2(\lambda,\tau) = \lambda S_1(\lambda) + S_2(\lambda,\tau),$

Using $(A_6)$ one can show that for $\lambda$ with re $(\lambda) \geq 0$, $c_1$ and $c_2$ are positive and bounded away from zero, Now since $A_{ij} > 0$ and $\tau \geq 0$ we find from the nature of the dependence of $g_{ij}$ on $\lambda$ that for all $\lambda$ $\mu \pm iw$ with $\mu \geq 0$ and $w \geq 0$ we have

(3.26) $$\left| F_1(\lambda,\tau) - P_2(\lambda,\tau) \right|$$

$$\geq \left| F_1(\lambda,0) \right| - F_2(\lambda,\tau) \Big|$$

$$\geq |\lambda|^2 - |\lambda|\{A_{11}+A_{22}+\| S_1(0) \|\} + \{A_{11}A_{22} - A_{12}A_{21} - \| S_2(0,\tau) \|\}$$

$$\geq \left[(|\lambda| - P)^2 + (q)\right]^2$$

Where

$p = \{A_{11}+A_{22}+\| S_1(0) \|\}/2,$

$q = \left\{4|(A_{11}A_{22} - A_{12}A_{21}) - \| S_2(0,T) \| | - |A_{11}+A_{22}+\| S_1(0) \| |^2\right\}/2,$

It will now follow, from (A) and (3.26), $F_1(\lambda,\tau) = F_2(\lambda,\tau)$ connot have roots with zero or positive real parts and this completes the proof.

In conclusion we remark that the condition (3.7) means that the net creation rates of each system is greater than utility; a condition of this type is well known in age wise classification of particles, atoms, molecules, cosmological phenomenon. (3.8) can be interpreted to mean that the intraspecific inhibition is higher than that of the interspecific interaction. The conditions of Theorem 4 are analytical and do not lend themselves for any worthwhile interpretation in terms of the model parameters. So also an interpretation can be recorded as a remark in respect of destruction rates of particles, properties, elements, or projections.

Note: The model is applicable to each and every module. As stated in the foregoing, there are possibilities of situation where an account itself be a receiver and it is also the one to which debit is raises. Such examples are legion in all the fields of finance, economics, political science and physics and chemistry, and of course mathematics. Classification based on time is the preeminently and predominantly suited premise for such model. We repeat that the model is to be read with each module We will expatiate the importance and implicational ramifications as we proceed with the work. Assortment-blend, collection-combinations, combo-gallimaufry, goulash-hash, hodgepodge-medley, mishmash-mixture, motley-pastiche patchwork-salmagundi, soup-stew, mixed bag-mélange, olio- olio podrida, assortment variety, Array choice, collection combination, combo diversity, hodgepodge jumble, kind medley, miscellany mishmash, mixed bag mixture, mélange potpourri, selection sort, brew concoction, beverage blend, broth compound, distillation drink, fermentation hash, hodgepodge infusion, instillation liquor, mélange miscellany, mishmash mixture, mélange potpourri, borscht bouillon, bowl brew, chowder concoction, consommé decoction, gumbo hodge-podge, olio porridge, potage potpourri, pottage puree, splash stock,

vichyssoise hodgepodge, mixture- mess, collection combination, goulash hash, jumble medley, miscellany mishmash, mixed bag mélange, olio patchwork, potpourri salmagundi, jumble clutter, disarrangement disarray, disorder farrago, gallimaufry garbage, goulash hash, litter medley, mess miscellany, mishmash mixture, muddle mélange, olio pastiche, patchwork potpourri, salmagundi scramble, shuffle snarl, tangle tumble, medley miscellany, assortment brew, collection combo, conglomeration farrago, hodgepodge jumble, melee mingling, mishmash mixed bag, mixture mélange, pasticcio pastiche, patchwork potpourri, salmagundi variety, mélange mixture, assortment combo, farrago gallimaufry, pasticcio pastiche, patchwork potpourri, salmagundi soup, stew miscellany, accumulation aggregation, anthology assortment, brew collectanea, combination combo, compilation conglomeration, accumulation diversity, farrago gallimaufry, melee mishmash, mix mixed bag, mixture muddle, mélange odds and ends, olio pasticcio, pastiche patchwork, potpourri salad, salmagundi smorgasbord, stew variety, mishmash jumble, medley mess, miscellany mixed bag, mixture mélange, olio patchwork, potpourri salmagundi, mixed bag array, assortment choice, collection combination, combo conglomerate, diversity hodgepodge, jumble medley, mélange miscellany, mishmash mixture, montage mélange, odds and ends, patchwork potpourri, smorgasbord sundries, mix/mixture assortment, combination, admixture adulteration, alloy amalgam, amalgamation assimilation, association batter, blend brew, combine combo, commixture composite, compound concoction, confection conglomeration, cross crossing dough fusion, goulash grab bag, hodgepodge hybrid, hybridization incorporation, infiltration interfusion, jumble mash, medley merger, mingling miscellany, mishmash mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, odds and ends assortment, bits and pieces, debris etcetera, hodgepodge jumble, knickknacks leavings, leftovers litter, medley melee, miscellaneous paraphernalia, miscellany mishmash, mixed bag motley, mélange notions, novelties oddments, olio particles, patchwork potpourri, rags refuse, remainder remains, remnants rest, rummage scraps, sundries sundry items, compilation mishmash, paste-up patchwork, potpourri reappropriation, reproduction synthesis, conglomeration cross, crossing dough, fusion goulash, interfusion jumble, mash medley, merger mingling, miscellany mishmash, mix mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, colluviums gallimaufry, hash mixture, pasticcio pastiche, conglomeration cross section, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multiformity multiplicity, mélange potpourri, range shift, variegation variousness, divergencies diverseness, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multifariousness multiformity, mélange potpourri, range shift, soup stew, variance variation, variegation variousness, oddment odds and ends, assortment bits, polymorphism assortment, change collection, combo conglomeration, cross section departure, discrepancy disparateness, divergencies diverseness, diversification diversity, fluctuation heterogeneity, heterogeneousness incongruity, intermixture many-sidedness, medley miscellaneousness, miscellany mishmash, mixed bag mixture, notions novelties, oddment oddments, olio particles, patchwork potpourri of the particles, interactions, accounts, transactions in the cosmos. The wisdom and understanding of our own divine essence is like a beaconing begorra, bureaucratic burdock, befuddling binary and digital digamma, that invites us inward toward ever deeper connection with all around us. Yet rather than look within, most in this world prefer to look outward for divine wisdom, thus sustaining a dependence upon a vast hierarchy that stretches between the individual and the Divine. In all our wanderings away from the Divine, humankind has obscured, obfuscated, obliterated, adjured, abnegated and abjured its most compelling features through a persistent belief systems in limitations arising out of the controlling mechanisms of this hierarchy of sub ordinates, superordinates and coordinates. Has my consciousness changed? Why that doubt? Twist and shake, ponderous novelty themes, pathos, poignancy in writing in different styles, accord-coetaneousness, coevality-coincidence, concurrence-conformity, conjunction-contemporaneousness, harmony, order, peace, simultaneousness- synchronicity exciting optimism of youth, power obmbudsman,pulsations and throbbing of pop, hard luck stories of people, modernity and mischief, grimalkins of

rurality, wobbling boats of lovers, saxophone in mouths all point to newer value orientations. The Divine dances and the “Grand design “is perpetrated outside of the confines of any hierarchical structures, be it feudal, organizational. It is complete within itself, and has a singular purpose of demonstrating the collective potential of all life within the universe. It is the archetype of perfection. It is the standard bearer of each soul's innate design and ultimate destiny. The essence of the Divine is far beyond mental conception, yet humanity's tendency is to resort to the limiting language of the hierarchical paradigm acceding- accession, accommodation- accord, accordance, acknowledging, adjustment-affiliation, affinity-alliance, amity-approving, arbitration-arrangement, assenting, authorizing, bargaining-compatibility, compliance--compromise, concert-concession, conformity-congruity, consistency-correspondence, endorsing, granting-harmony in conformity to pen chance ,predilections of the protagonist to define and understand it. Searching for Connection and Wholeness calls for not lumping elements like me made of clay feet and weak knees, owing allegiance to the fringe, without courante and courage of conviction and will for vindication, visual restlessness and exaggerated narratives of self to its own, exciting potentialities available, terrestrial passion, and skew of scams all bear an infallible demonstration for new order of consciousness. Theorists have posited a virtual smorgasbord of objects to account for dark matter, although many of them have fallen prey to observational constraints. As leading possible candidates for baryonic dark matter, there are black holes (large and small), brown dwarfs (stars too cold and faint to radiate), sun-size MACHOs, cold gas, dark galaxies and dark clusters, to name only a few. The range of particles that could constitute nonbaryonic dark matter is limited only slightly by theorists' imaginations. The particles include photinos, neutrinos, gravitinos, axions and magnetic monopoles, among many others. Of these, researchers have detected only neutrinos--and whether neutrinos have any mass remains unknown. Experiments are under way to detect other exotic particles. If they exist, and if one has a mass in the correct range, then that particle might pervade the universe and constitute dark matter. But these are very large "ifs." To a great extent, the details of the evolution of galaxies and clusters depend on properties of dark matter. Without knowing those properties, it is difficult to explain how galaxies evolved into the structures observed today. As knowledge of the early universe deepens, I remain optimistic that we will soon know much more about both galaxy formation and dark matter. What we fail to see with our eyes, or detectors, we can occasionally see with our minds, aided by computer graphics. Computers now play a key role in the search for dark matter. Historically, astronomers have focused on observations; now the field has evolved into an experimental science. Today's astronomical experimenters sit neither at lab benches nor at telescopes but at computer terminals. They scrutinize cosmic simulations in which tens of thousands of points, representing stars, gas and dark matter; interact gravitationally over a galaxy's lifetime. A cosmologist can tweak a simulation by adjusting the parameters of dark matter and then watch what happens as virtual galaxies evolve in isolation or in a more realistic, crowded universe. Old consciousness will not do. Newer one has to be imbibed and cultivated. When people are unaware of their inner wisdom and wholeness, they tend to search for order and security outside themselves. That is my personal experience. There is always search for sanctity of emotion, aesthetic expression, pursuit of harmony, unraveling realms, syncopated ending and stillness of mind and shiver of heart at the mirror broken somewhere are the strains of newer wisdom if it can be called one. Same questions shall be asked and same shall be answer in the similar vein. Uncertain of their place within the hierarchical order of the world, they define themselves based upon their insecurities. Individuals thus become only pieces of their wholeness and like shards of glass from a beautiful vase; they bear little resemblance to their aggregate beauty. Within the hierarchy, many in high positions of power have taken advantage of our collective insecurities in attempting to guide the development of all humankind. They have obscured the direct connection between the individual and the Divine through a variety of means designed to intercede between our inner essence and our divine source. Each individual can come to know them to be free of all forms of institutional and organisational and societal hierarchical control. This is not to imply that we should not trust others or join together in bonds of friendship and community. It is simply a reminder that relative truth is constantly shifting in the hands of those who desire to control. And even when the motive for their controlling behaviors may be of good will, it is still a form of control. When the revealers of "truth" within the hierarchy withhold and suppress information, they are usually positioning themselves to maintain or acquire more power

rather than to disseminate empowerment to all. The desire for connection and wholeness is a fuel that drives us to seek out and explore the hierarchical paradigm. This inner longing provides us with the motivation to seek help and guidance from a specific group within the hierarchy, and in so doing, cultivate a sense of belonging and connection. Furthermore, the hierarchical paradigm is a stage whereby we develop a sense of connection to some grand, encompassing vision. This is why the hierarchy nurtures prophets who point toward a greater vision. Spiritual leaders are able to peer deeply beneath the surface reality of life and experience how intricately connected every life form is, and how the composite of all life is intelligent far beyond measure. These visionary leaders can thus interpret reality through their personal abilities to perceive and express life's dimensional depth and limitless intelligence. Yet no one is able to articulate life's full dimensional depth and breadth with the tools of language. They can only, at best, describe their interpretation or their impressions. In actuality, all of us are able to peer beneath the surface reality of life and perceive a unique vision of the universe. We require only time and intention to develop our own interpretations. And this is precisely what many great spiritual leaders have taught. Life's deeper essence is not an absolute to be experienced by the chosen few, but an evolving, dynamic intelligence that wears as many faces as there are life forms. No group or people have the exclusive portal into the universe by which the Divine expresses itself in all its majesty. The portal is open and available to all, because the Divine is within all things. Those recognized as great prophets each produced a vision of the universe beyond what was currently defined by the hierarchy. Because their interpretations were articulated with authority and depth of insight, they became a target of debate among various groups in the hierarchy. This debate then created a polarity of belief. A sympathetic constituency emerged to defend and embellish their leader's interpretation, while established groups held it in contempt of previously held beliefs. Invariably, the leader's vision became confined and shaped into dogma by followers who desired to create a new religion or sect. Thus, this infusion of fresh insight quietly receded into the hands of the hierarchy, where its deeper meaning was obscured by the very fact that it was incorporated into a rigid structure that both protected and promoted it. A new paradigm is emerging that promotes a clear connection of individual consciousness to the compelling features of the Divine without the intervention of a hierarchy. This is when the fables and myths of history step into the light and become known as they were originally intended. This is the time when language will be transformed into a new form of communication that breaks down all barriers of control. Personal transformation, through the awakening of inner wisdom and understanding, is the pathway into wholeness. The transformation paradigm is initiated simply by the recognition that rather than the dependency-inducing ways of the hierarchical paradigm, there are accelerated, independent pathways that bypass the hierarchy and lead to self-mastery. These new pathways lead to the divine wisdom and understanding that is present within all of us. This wisdom can be accessed through the practice of three principles of transformation: seeing the Divine in all, nurturance of life, and gratitude. The application of these life principles disengages individuals from the controlling elements of the hierarchy, thus initiating the transformation experience. As there are relative truths, there are relative freedoms. As individuals evolve through the hierarchical process, an ever-increasing sense of freedom is gained, yet external forces continue to exert control through limiting language, confining belief systems, and more. These controlling influences lead to anatomical architecture, body, build, cagey carcass, construction, enclosure, fabric, flounce, form, framework, fringe, groundwork, hem, mount, mounting outline, physique, scaffold, scaffolding scheme, setting shell, stage, structure, support, system, trim, trimming, truss- valance continued reliance upon the hierarchy as it unceasingly attempts to impose a sense of inequality between us. The underlying equation of the hierarchical paradigm is: individual + hierarchy = God connection. In the case of personal transformation it is: individual + inner wisdom and understanding divine equality with all. The time has come to integrate the dominant paradigm of the hierarchy with the liberating transformation paradigm. That shall redeem all promises made to oneself. There shall be a Matrix force or balancing force that shall remain UN satisfies until conservativeness is reached. This integration occurs naturally once we have fully explored the two paradigms and develop a synthesis paradigm whereby transformation is attained by searching for connection and wholeness through our own inner wisdom and understanding. It is this combination of release from dependency on the hierarchy and transformation from within that initiates the synthesis paradigm. Once we take responsibility for our own transformation and

integration, it does not mean that the hierarchy is to be shunned or avoided. The hierarchy is quite benign as a manipulative force and merely represents one important stage in the journey toward wholeness. What is being set in motion now is the initial preparation for these paradigm shifts. More specifically, these paradigms will be simultaneously played out over the coming years. As always, it will be the choice of each one of us as to which paradigm we embrace in our journey. All of the highest imaginings of humanity are yet unaware of our deepest foundation. We have sought the upper reaches of the building, yet remain unaware of the foundation's design. It is here, at the very core of existence that the Divine is bursting forth with its creative energy while simultaneously reintegrating with its invitation to wholeness. It is here that equality is realized, not in the lofty places of relative truth lodged in the hierarchy, but rather in the deepest part of the foundational plan of life's original source and ultimate destiny. The origin and destiny of existence is the tone of equality in life. Listen for this tone – this frequency of vibration – and follow it back into the very foundation from whence all things arise and return. Consider these words as symbols only. Feel the truth that stands behind these words, and tap into this empowering energy force that reaches out for you. Know it as a tone or vibration – a resonance that waits for you around every corner in which your life will turn. It is a beacon of the Divine gathering itself into the form of language in order to usher you to a place from which you can experience the formless tone of equality – the bypass of limitation. It is the primal language of our divine source that bestows to you the freedom to generate your own deepest beauty in the expression of your highest truth.(Reference: Wikipedia)

## Gopalasamy Model For Procrastinated Conducive, Felicitous Counter Sallies In Two Specified Systems

Abstract

Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of decay rate of perturbations. Sole intention, main objective and primary concern of the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are

counted in the same vein as that of predator prey. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 13. Preliminary remarks

Almost all the equations of population ecology have been derived starting from the following simple format (Note this can be used for the development of the creed and type of systemic types in case of cosmology and astrophysics and fiancé and econometrics. To give an example of finance, we can think of a account of the company giving rise to various accounts of its employees, and that is the contribution of the single account for the development of the total system types of the given configuration). It may be account, particle, atom, molecule, proposition, problem or any such thing. We shall delineate and disseminate information and particulars as and when the model shall be applied to the problem in question.

$$\frac{dN(t)}{dt} = \left\{ \begin{array}{l} \textit{an individual's contribution} \\ \textit{to population growth} \end{array} \right\} N(t)$$

(1.1)

where $N(t)$ denotes the size of a population of a single species (or biomass) with overlapping generations at time $t$, and $dN(t)/dt$ is the rate of change of the population size at time $t$. Subsequently one makes an assumption regarding the factor in the brackets in (1.1); in particular if one assumes that an individual entity's contribution to the population growth is denoted by a function say, f(t, N), defined suitably for all t>0, N0 then one obtains from (1.1) the so called Kolmogorov formulation in the form

$$\frac{dN(t)}{dt} = f(t, N(t))\, N(t).$$

(1.2)

Various choices of $f$ together with some ecologically plausible assumptions such as the temporal constancy of the environment and density dependent effects in $f$ lead to various equations. For instance if $f(t, N) \equiv r$ (a positive constant) one obtains the Malthusian formulation in the form

$$\frac{dN(t)}{dt} = rN(t)$$

And if one assumes $f(t, N) \equiv r - (r/K)N$ for some positive constants r and K one gets the familiar logistic equation of population biology

$$\frac{dN(t)}{dt} = rN(t)\left\{1 - \frac{N(t)}{K}\right\}$$

In order to device models with time delayed effects, such effects are introduced in the average growth rate $F(t, N)$ of (1.2); for instance if $f(t, N) \equiv r - (r/K)N(t - \tau)$ for some constant $\tau > 0$ we obtain Hutchinson's[19]delay loistic equation

$$\left\{\frac{dN(t)}{dt} = rN(t)\right\}\left\{1 - \frac{N(t-\tau)}{K}\right\}. \quad (1.3)$$

If for instance

$$F(t, N) = r - (r/K)N(t) - \int_{-\infty}^{t} k(t - s)F(N(s))\,ds$$

Where $k$ is a piecewise continuous nonnegative valued function such that $\int_0^\infty k(s)\,ds = 1$ with some appropriate function $F$, we will get Volterra's model of a population which pollutes its environment and pollution has an accumulative and toxic effect; such an equation is of the form

$$\left\{\frac{dN(t)}{dt} = N(t) = \right\}\left\{r - (r/K)N(t) - \int_{-\infty}^{t} k(t - s)F(N(s))\,ds\right\} \quad (1.4)$$

Which has been studied in different formats by several authors (Cushing [12], Cohen *et al.* [10], and Gopalasamy [16])

One of the significant advantages of deriving various models from the prototype (1.2) is that if the initial; conditions (initial population sizes) are nonnegative then the nonnegativity of the population size $N(t)$ for $t > 0$ follows from the fact that solutions of (1.2) satisfy

$$N(t) = N(0) \exp\left[\int_0^1 f\left(s, N(s)\right)ds\right]$$

Which implies that $N(t) \neq 0$ if $N(0) \neq 0$.

While the time delayed logistic equation (1.3) explains some of the observed oscillations of certain populations, its derivation has not been entirely satisfactory. For instance it has been pointed out by Ricklefs ([27], page 488) that "time lags might be expected to occur primarily in stabilizing processes involving production rather than death because death is an immediate response to environmental change". Recently Cushing [13] has examined the effects of response and maturation delays in a number of single and multi system models; as it has been noted by Cushing [13], the models involving delayed reproduction or not necessarily derivable from the Kolmogorov type systems of the form (1.2) by suitably defining the average growth rate $(1/N)\,(dN/dt) = f(t, N(t))$. While Cushing [13] has been mostly concerned with the possibility of delay induced oscillations we will consider a class of delays in reproduction and interspecific interaction which maintain the stability of the systems.

Our starting point to drive models with delayed reproduction is the following balance equation which assumes that there is no immigration or emigration. Account, particle, atom, molecule, proposition, problem et al

$$\frac{dN(t)}{dt} = \text{birth rate} - \text{death rate.}$$

(1.5)

Now for instance if we consider a population of adult flies then the birth rate at time $t$ depends on the adult population at time $t - \tau$ where $\tau$ is the time required for the larvae to become adult flies. If the birth and death rates are governed by density dependent factors then we have from (1.5) the following time delayed model

$$\frac{dN(t)}{dt} = b\left(N(t-\tau)\right) - m(N(t))$$

(1.6)

Where the functions $b$ and $m$ denote the density dependent birth and death rates of Account, particle, atom, molecule, proposition, problem et al respectively. If the time delay is continuously distributed then one can consider in the place of (1.6) an equation of the form

$$\frac{dN(t)}{dt} = b\int_{-\infty}^{t} k\left(t-s\right)N(s)ds - m\left(N(t)\right)$$

Or equivalently

$$\left\{ \frac{dN(t)}{dt} = b \int_0^{\infty} k(s) N(t-s) ds - m(N(t)) \right. \quad (1.7)$$

Where the delay kernel *k* denotes the weight of the delayed of the delayed effects.

It is not always obvious that suitable nonnegative initial conditions for (1.6) and (1.7) will imply the nonnegativity of the solutions of (1.6) and (1.7). The purpose of this article is to show that in two species systems of competitive and prey- predator communities, delayed reproduction and delayed interactions donot destabilise the systems if the self- regulating negative feedback effects are sufficiently strong compared to interspecific interactive effects and if the self regulating effects are realised with no time delays. When we say that time delays do not destabilise, we will mean that delays do not render an asymptotically (locally) stable steady state unstable in such a manner that the loss of stability either leads

To a delay – induced bifurcation to persistent and undamped oscillations (as in the case of Hopf- bifurcation) or the relevant linear variational system has unbounded solutions; that is a mathematical instability is not induced by delays in the sense, that the roots of the characteristic equation are pushed to the imaginary axis or to the right of it on the complex plane. However delays can make an otherwise stable system less stable which means that the rate of convergence to zero of solutions of the variational system can decrease with an increase of delay and the variational system can even oscillate (perhaps violently) before converging to the trivial state. Such a reduction in stability can happen even when all the roots of the characteristic equation have negative real parts. For our purposes such a reduced stability is still stability. In order to see the effects of delays on such less stable systems we will propose below a method of estimating the rate of decay of local perturbations based on the variational system and indirectly relying on the characteristic equation. For a special class of two species models with a single delay we have obtained (Gopalsamy [16]) estimates of decay rates (or return times) usingthe characteristic equation; such a procedure is almost impossible in the case of many species systems with several delays. The method we propose in Section 6 below for the estimate of decay rates can be adapted to *n*- species systems with $n^2$ discrete delays.

Let $x_1(t)$ and $x_2(t)$ denote the population sizes (or biomasses) of two systems competing for a common pool of resources with common pool of resources.. Let $b_i$ and $m_i$ $(i = 1, 2)$ denote the respective density dependent birth (potential) and death rates so that in the absence of any time delays the population sizes are governed by (of

Account, particle, atom, molecule, proposition, problem et al). In addition as usual, we assume that the system studied follow Haimovici's equations, for which the dove tailing explanation has been appended

$$\frac{dx_1(t)}{dt} = b_1\left(x_1(t)\right) - m_1(x_1(t)), x_2(t))$$

$$\frac{dx_2(t)}{dt} = b_2\left(x_2(t)\right) - m_2(x_1(t)), x_2(t))$$

(2.1)

In order to make the system (2.1) represent a competition system we will make the following assumptions on the birth and death rates; $_1$, $m_1$ ($i$ = 1, 2) are continuous with continuous partial derivatives for all $x_1$ 0 ($I$ = 1, 2); also we assume (see for instance Rescigno and Richardson [26]) $\frac{\partial b_i}{\partial x_1} > 0, \frac{\partial m_1}{\partial x_j} > 0$ for $x_i > 0$ ($i, j$ = 1, 2), (2.2)

(xix) $b_1(0) = 0$, $m_1(0, x_2) \equiv 0, b_2(0) = 0$, $m_2(x_1, 0) \equiv 0$,

(2.3)For some $x^*_{1,} > 0, x^*_{2,} > 0$ we have

$$b_1(x^*_1, -m_1(x^*_1, 0) = 0,$$

$$b_2(x^*_2, -m_1(0, x^*_2) = 0,$$

(2.4)

(xx) There exist positive constants $\delta_1$, $\delta_2$ such that

$$b_1(\delta_1) - m_1(\delta_1, x_2) < 0,$$

$$b_2(\delta_2) - m_2(x_1, \delta_2) < 0,$$ for $x_1$ 0, $x_2$ 0,

(2.5)

(xxi) There exists a point $(\alpha, \beta), \alpha > 0, \beta > 0$ such that

$$b_1(\alpha) - m_1(\alpha, \beta) = 0 = b_2(\beta) - m_2(\alpha, \beta).$$

(2.6)

The conditions on $b_i$ in (2.2) mean that the potential (not average) birth rates are density dependent and any negative density dependent feedbacks to the opened or birth rate of Account, particle, atom, molecule, proposition, problem et al rates are included in the destruction rates; the conditions on $m_i$ indicate intraspecific and interspecific confederation needs and necessities. (2.3) implies that (0, 0) is a (trivial) steady state of the system (2.1) as it is in

the case of all ecological models; (2.4) means that in the absence of either of the two species the other one has a nontrivial steady state; (2.5) implies that each of the two populations cannot grow unboundedly since when $x_1(t) > \delta_i, \frac{dx_1}{dt} < 0$. (2.6) means that there exists a nontrivial steady state $(\alpha, \beta)$ in the interior of the positive quadrant of the $(x_1, x_2)$- plane. These assumptions are standard and are minimal in the opinion of the author.

An immediate example of (2.1) –(2.6) is the familiar Volterra –Lotka competition system described by

$$\left.\begin{aligned}\frac{dx_1(t)}{dt} &= r_1x_1(t) - a_{11}x_1^2(t) - a_{12}x_1(t)x_2(t),\\ \frac{dx_2(t)}{dt} &= r_2x_2(t) - a_{21}x_1(t)x_2(t) - a_{22}x_2^2(t),\end{aligned}\right\} \quad (2.7)$$

Where $r$, $a_{ij}$ , $(i, j = 1,2)$ are positive constants satisfying the condition

$$\frac{a_{11}}{a_{21}} > \frac{r_1}{r_2} > \frac{a_{12}}{a_{22}}.$$

One of the first questions to be resolved for (2.1) is the following: if $x_i(0) > 0$, $x_2(0) > 0$ does it follow that the corresponding solutions of (2.1) satisfy the condition $x_1(t) > 0$, $x_2(t) > 0$ for those $t$ for which a solution of (2.1) exists? In the $(x_1, x_2)$ phase plane of (2.1) we have $(0, 0), (x_1^*, 0), (0, x_2^*), (\alpha, \beta)$ as the steady states of (2.1) and also it will follow from (2.3) that the co ordinate axes

$x_1 = 0$ and $x_2 = 0$ are invariant sets f (2.1) ; such invariance of the coordinate axes will imply that if $x_1(0) > 0$, $x_2(0) > 0$ then the corresponding population trajectory $(x_1(t), x_2(t))$ of (2.1) cannot reach the outside of the nonnegative quadrant $x_1 \geq 0$, $x_2 \geq 0$.

Our plan is to derive first a set of sufficient conditions for the coexistence of the two competing species and then in the next section show that under the same conditions the two species can coexist if there are time delays in the reproduction and interspecific interactions. We recall that the two competing species described by (2.1) are said to coexist if and only if the nontrivial steady state $(\alpha, \beta)$ is (at least locally) asymptotically stable. To examine such a stability of $(\alpha, \beta)$ we derive the corresponding linear variational system for the perturbations $X_1, X_2$ where

$$x_1(t) \equiv \alpha + X_1(t), \qquad x_2(t) \equiv \beta + X_2(t)$$

Such a variational system is found to be

$$\left.\begin{aligned}\frac{dX_1(t)}{dt} &= \beta_{11}X_1(t) + \beta_{12}X_2(t),\\ \frac{dX_2(t)}{dt} &= \beta_{21}X_1(t) + \beta_{22}X_2(t),\end{aligned}\right\} \quad (2.8)$$

Where

$$\beta_{1J} = \left(\frac{\partial b_1}{\partial x_J} - \frac{\partial m_1}{\partial x_j}\right) \text{ evaluated at } (\alpha, \beta) \qquad (i, j = 1,2). \quad (2.9)$$

The steady state $(\alpha, \beta)$ of (2.1) is locally asymptotically stable if and only if the trivial solution (0, 0) of (2.8) is asymptotically stable. We can derive the following result from elementary considerations.

In the system (2.1)-(2.6) assume the following:

$$\left.\begin{aligned}\frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1},\\ \frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1},\end{aligned}\right\} \text{ at } (\alpha, \beta) \quad (2.10)$$

Then the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (2.1) – (2.6).

The proof follows easily if we note that the characteristic equation associated with the linear system (2.8) is given by

$$\text{Det}\begin{pmatrix}\lambda-\beta_{11} & -\beta_{12}\\ -\beta_{21} & \lambda-\beta_{22}\end{pmatrix} = 0$$

Or equivalently,

$$\lambda^2 - \lambda(\beta_{11} + \beta_{22}) + \beta_{11}\beta_{22} - \beta_{12}\beta_{21} = 0. \quad (2.11)$$

It follows from (2.2) and (2.10) that,

$$\beta_{11} < 0, \beta_{11}\beta_{22} - \beta_{12}\beta_{21} > 0$$

And hence the roots of (2.11) have negative real parts implying the asymptotic stability of the trivial solution of (2.8)

If apply the condition (2.10) to the Volterra-Lotka model (2.7) we note that when a nontrivial steady state say $(\alpha, \beta)$ exists for (2.7), (2.10) leads to

$a_{11}\alpha > a_{21}\,\beta \; and \; a_{22}\,\beta > a_{12}\,\alpha,$

Which together imply$a_{11}\,a_{22} > a_{12}a_{21}$. If $(\alpha,\beta)$ exists it is known that $a_{11}\,a_{22} > a_{12}a_{21} > 0$ is a sufficient condition for the (local) asymptotic stability of $(\alpha,\beta)$.

The condition (2.10) which guarantees the coexistence of competing populations in (2.1)-(2.6) simply means that the intraspecific self – regulating feedback effects $\partial m_i/\partial x_1\,(i=1,2)$ bare higher than its own positive feedback $\partial b_1/\partial x_1\,(i=1,2)$ together with its negative (competitive) effects on its competitor $\partial m_I/\partial x_J\;(i,\;j=1,2;i\neq j\,)$.

## 14. Coexistence under delayed production and interactions, transactions

We will now establish that two competing species with a sufficiently strong intraspecific regulation can coexist in the presence of delayed reproduction and delayed interspecific interactions. It is well known that if there are delays in the self – regulating intraspecific effects, stationary coexistence may not always be possible (Gopalsamy [16], Gopalsamy and aggarwala [15], Shibata and Saito [28]). Delayed reproduction is more general and widespread in most biological populations. To build a competition model with such delays we proceed as follows.

Let $\tau_{iJ}\;(i,j=1,2)$ be a set of nonnegative constants with $\tau = \max\,\tau_{ij}\,,i,j=1,2,$ and suppose that the two ompeting species display delayed reproduction and delayed interspecific interactions with no delays in the self – regulating negative feedback effects. Such a competition system in a constant environment can be modeled by the autonomous delay differential system

$$\frac{dx_1(t)}{dt} = b_1\big(x_1(t-\tau_{11})\big) - m_1\big(x_1(t),x_2(t-\tau_{12})\big),$$

$$\frac{dx_2(t)}{dt} = b_2\big(x_2(t-\tau_{22})\big) - m\big(x_1(t-\tau_{21}),x_2(t)\big), \qquad (3.1)$$

Where both the creation and destruction rates and the morality rates $b_1$ and $m_1\;(i=1,2)$ satisfy the same conditions (2.2)-(2.6) as in the previous section. Along with (3.1) we will

Suppose that the initial population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) sizes are provided in the form

$$x_i(s) = \phi_i(s)\;\;0, \qquad s\in[-\tau,0], \qquad i=1,2. \qquad (3.2)$$

(Note that the specification of $\phi_i$ need not be on the entire interval$[-\tau,0]$; but such a hypothesis is convenient.) Since (3.1) is not of the standard Kolmogorov-type, we have to establish that the solutions of (3.1) will remain nonnegative for nonnegative initial conditions.

To prove the nonnegativity of solutions of (3.1)-(3.2), let $\left\{\psi_1^{(n)}(s)\right\}, s\in[-\tau,0], i=1,2; (n=1,2,\ldots)$ be a sequence of strictly positive continuous functions such that

$$\lim_{n\to\infty}\psi_1^{(n)}(s), = \psi_1\;(s), \qquad s\in[-\tau,0], \qquad i=1,2 \qquad (3.3)$$

Let $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ be the solution of (3.1) corresponding to the initial condition

$$\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \left\{\psi_1^{(n)}(s), \psi_2^{(n)}(s)\right\}, \quad s \in [-\tau, 0], \tag{3.4}$$

Consider the solution $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ for $t \in [0, \tau]$ (assuming it exists on the whole of $[0, \tau]$). Suppose $x_1^{(n)}$ does not remain nonnegative for all $[0, \tau^*]$; then there exists a $t^* = [0, \tau^*](\tau^* =$ the positive minimum of $\tau_{ij}$ $(i, j = 1, 2))$ for which

$$x_1^{(n)}(t^*) = 0 \ and \ x_1^{(n)}(t)(t^* - \tau_{11})) > 0.$$

It will follow from (3.1) and the positivity of the initial condition that

$$\frac{dx_1^{(n)}(t^*)}{dt} = b_2\left(x_1^{(n)}(t^* - \tau_{11})\right) > 0 \tag{3.5}$$

Where we have used the properties of $b_i$ and $m_i$ from (2.2)-(2.6); it is found that (3.5) contradicts the definition of $t^*$. Thus $x_1^{(n)}(t) > 0$ on $[0, \tau^*]$; similarly $x_2^{(n)}(t) > 0 \ on \ [0, \tau^*]$. We can repeat this procedure for intervals of the form $[\tau^*, 2\tau^*], [2\tau^*, 3\tau^*] \ etc$. Thus it will follow that so long as $(x_1^{(n)}(t), x_2^{(n)}(t))$ exists, we have

$$x_1^{(n)}(t) \geq 0 \ x_2^{(n)}(t) \geq 0.$$

If we now consider the limits as $n \to \infty$, we get

$$\lim_{n\to\infty}\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \{x_1(t), x_2(t)\} \tag{3.2}$$

With $x_1(t)$ ,, $x_2(t)$ where $(x_1(t), x_2(t))$ is the solution of (3.1)-(3.2) and this is a consequence of the continuous dependence of solutions of (3.1) on initial conditions (Hale [17], page 41).

A second question for (3.1) with (2.2)-(2.6) is whether solutions of (3.1) exist for all $t > 0$? suppose now that solutions of (3.1) do not exist for all $t$ ; then there exists a $t_1 > 0$ such that at least for one of either of $x_1 \ or \ x_2$ we have

$$\lim_{t\to t_1-} x_i(t) = \infty. \tag{3.7}$$

To be specific let us suppose that $\lim_{t\to t_1-} x_i(t) = \infty$; let $t_2$ be the first instant of time for which $x_1(t_2) = \delta_1, t_2 < t_1$. It will follow from (3.1) that

$$\frac{dx_1(t_2)}{dt} = b_1(x_1(t_2 - \tau_{11})) - m_1(x_1(t_2), x_2(t_2 - \tau_{12}))$$

$$< b_1(x_1(t_2)) - m_1(x_1(t_2), x_2(t_2 - \tau_{12}))$$

$$< b_1(\delta_1) - m_1(\delta_1, x(t_2 - \tau_{12})) < 0 \quad (3.8)$$

Which contradicts the definition of $t_2$. Thus solutions of (3.1), (2.2) – (2.6) exist for all $t \geq 0$.

We are now ready to consider the effects of time delays in reproduction and interspecific interactions on the coexistence of the two competing species. The steady state $(\alpha, \beta)$ of (2.1) is again a steady state of (3.1). To examine the (local) asymptotic stability of $(\alpha, \beta)$ in (3.1) we let

$$x_1(t) = \alpha + Y_1(t), \quad x_2(t) = \beta + Y_2(t)$$

In (3.1) and derive that the linear variational system governing the perturbations $Y_1, Y_2$ is given by

$$\frac{dY_1(t)}{dt} = -m_{11}Y_1(t) + b_{11}Y_1(t - \tau_{11}) - m_{12}Y_2(t - \tau_{12}),$$

$$\frac{dY_2(t)}{dt} = -m_{22}Y_2(t) + b_{21}Y_1(t - \tau_{21}) - m_{22}Y_2(t - \tau_{22}),$$

Where

$$m_{ij} = \frac{\partial m_1}{\partial x_j}, \quad b_{ii} = \frac{\partial b_i}{\partial x_i} \; at \; (\alpha, \beta). \quad (3.10)$$

It is well known (Bellman and Cooke [3], page 336) that the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) if the trivial solution of the variational system (3.9) is asymptotically stable. We formulate our first result in the following form whose proof is given in the appendix.

*THEOREM 3.1. Assume that in the two systemic types competition model system (3.1) with delayed density dependent birth and death rates $b_i$ and $m_i$ $(i = 1, 2)$ satisfy the conditions (2.2)-(2.6). Suppose that the conditions (2.10) hold. Then the nontrivial steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) whatever the delays be $\tau_{ij} \geq 0$ $(i = 1, 2)$.*

Let $y_1(t) and\ y_2(t)$ respectively denote the population sizes (or biomasses) of a prey species and a predator species which feeds exclusively on the single prey species in a temporally constant environment. In the format of (1.5) the following system will describe the rates of change of the two population sizes:

$$\frac{dy_1(t)}{dt} = f_1(y_1(t)) - g_1(y_1(t), y_2(t)),$$

$$\frac{dy_2(t)}{dt} = f_2(y_1(t)) - y_2(t) - g_2(y_2(t)), \quad (4.1)$$

Where $f_i\ and\ g_i\ i = 1,2)$ respectively denote the birth and death rates of the two species; in order to make (4.1) denote the dynamics of a prey – predator species we will assume that $f_i\ and\ g_i$ satisfy the following conditions (see for instance Rescigno and Richardson [26]); $f_i\ and\ g_i$ are continuous with continuous partial derivatives for all $y_i \geq 0$ $(i = 1, 2); also\ we\ require$

$$\frac{\partial f_1}{\partial y_1} > 0; \quad \frac{\partial g_1}{\partial y_1} \; 0; \quad \frac{\partial g_1}{\partial y_2} \; 0,$$

$$\frac{\partial f_2}{\partial y_1} \; 0; \quad \frac{\partial g_2}{\partial y_2} \; 0; \quad \frac{\partial g_2}{\partial y_2} \; 0, \quad \text{Where } y_1 \; 0, y_1 \; 0, \tag{4.2}$$

(xxv)

$$f_1(0) = 0, \quad g_1(0, y_2 \equiv 0,$$

$$f_2(0, y_2) = 0 = f_2(y_1, 0); \quad g_2(0) = 0, \tag{4.3}$$

(xxvi) For some constant $y_1^* > 0$ we have

$$f_1(y_1^*) - g_1(y_1^* 0) = 0, \tag{4.4}$$

(xxvii) There exists a point $(\alpha^* \beta^*), \alpha^* > 0, \beta^* > 0$ such that

$$f_1(\alpha^*) - g_1(\alpha^*, \beta^*) = 0,$$

$$f_2(\alpha^*, \beta^*) - g_2(\beta^*) = 0, \tag{4.5}$$

(xxviii) There exist positive constants $n_1, n_2$ such that

$$f_1(\eta_1) - g_1(\eta_1, y_2) < 0, y_2 \; 0,$$

$$f_2(\eta_1, \eta_2) - g_2(\eta_2) < 0 \tag{4.6}$$

The conditions (4.2) mean that the system (4.1) is of the oppugnacy systems type; (4.3) will imply that (0,0) is a trivial steady state of (4.1); (4.4) means that in the absence of the predator, the prey species has a nontrivial (positive) steady state;

(4.5) implies that (4.1) has a nontrivial steady state in the interior of the positive quadrant of the state species of (4.1); (4.6) can be shown (see below) to imply that no credit type and debit type system types can grow unboundedly.

An example of (4.1)-(4.6) is the familiar Volterra-Lotka prey-predator system

$$\frac{dy_l(t)}{dt} = r_l y_l(t) - [a_{ll} y_l^2(t) + a_{12} y_l(t) y_2(t)],$$

$$\frac{dy_2(t)}{dt} = a_{12} y_1(t) y_2(t) - [r_2 y_2(t) + a_{22} y_2^2(t)],$$

Where $r_l, a_{ij}\,(i, j = 1,2)$ are positive costants such that

$$a_{21} r_1 - r_2 a_{11} > 0.$$

As in the previous case of competition we have to verify that when $y_1(0) > 0, y_2(0) > 0$ the corresponding solution of (4.1) will be nonnegative for $t \; 0$ for which the solution exist. Since the co-ordinate axes $y_1 = 0, y_2 = 0$ are invariant sets for (401) by virtue of (4.3), it will follow that the solution of (4.1) starting from the interior of the positive quadrant cannot reach the outside of that quadrant. This is a consequence of the invariance of the co-ordinate axes for (4.1)

To prove that solutions of (4.1)-(4.6) exist for all t 0 let us suppose that $y_l 0 \; and \; y_2(0) satisfy \; 0 < y_l(0)$ satisfy $0 < y_l(0) < \eta_l, 0 < y_2(0) < \eta_2$ ( on the other hand if either one of $y_l(0) \; \eta_l and \; y_2(0) \; \eta_l$ or both hold, it will follow that the corresponding $y_l(t)$ will decrese by virtue of (4.6) at least for small $t \; 0,$ ). Suppose $y_l(t)$ does not

exist for all $t$ ; then for some $t_l < \infty\ lim_{t\to t_l-} y_l(t) = \infty$and let $t_2 > 0$ be the first instant for which $y_1(t_2) = \eta_1$; then we have from (4.1) and (4.6) that

$$\frac{dy_1(t_2)}{dt} = f_1(\eta_1) - g_1(\eta_1, y_2) < 0\ (assumimg\ y_1(0) < \eta_1)$$

Which contradicts the definition of $t_2$ showing that $y_1(t)$ exists for all $t$ ; if $y_2(t)$ is not defined for all $t$ , then for some $t_3$ we will have $lim_{t-t_3-} y_2(t) = \infty$;now let $t_4$ be the first instant for which $y_2(t_4) = \eta_2$; then from the properties of $f_2$ we have

$$\frac{dy_2(t_4)}{dt} = f_2(y_1(t_4), \eta_2) - g_2(\eta_2)$$ $$< f_2(\eta_1, \eta_2) - g_2(\eta_2)$$

$$< 0 (by\ (4.6) assuming\ 0 < y_1(0) < \eta_1, 0 < y_2(0) < \eta_2)$$

Which again contradicts the definition of $t_4$. Thus $y_2(t)$ is defined for all $t$ .

We will now derive a set of sufficient condition for the (local) asymptotic stability of the steady state $(\alpha^*, \beta^*)$ of (4.1). in the next section we will show that under the same conditions and in the presence of various delays , the steady state $(\alpha^*, \beta^*)$remains (locally) asymptotically stable.

To examine the asymptotic stability of the nontrivial steady state $(\alpha^*, \beta^*)$ in (4.1) we let

$$y_1(t) = \alpha^* + U_1(t), \qquad y_2(t) = \beta^* + U_2(t)$$

In (4.1) and derive the liner variational system in the perturbations $U_1, U_2$ as follows:

$$\left.\begin{aligned} \frac{dU_1(t)}{dt} &= \frac{\partial f_1}{\partial y_1} U_1(t) - \frac{\partial g_1}{\partial y_1} U_1(t) - \frac{\partial g_1}{\partial y_2} U_2(t), \\ \frac{dU_2(t)}{dt} &= \frac{\partial f_2}{\partial y_1} U_1(t) - \frac{\partial f_2}{\partial y_2} U_2(t) - \frac{\partial g_2}{\partial y_2} U_2(t), \end{aligned}\right\} \quad (4.7)$$

Where all the partial derivatives in (4.7) are evaluated at $(\alpha^*, \beta^*)$. The steady state $(\alpha^*, \beta^*)$ of (4.1) is (locally) asymptotically stable if and only if the trivial solution of (4.7) is asymptotically stable which will be the case if and only if the roots of the characteristic equation

$$\text{Det} \begin{pmatrix} \lambda - (f_{11} - g_{11}) & g_{12} \\ -f_{21} & \lambda - (f_{22} - g_{22}) \end{pmatrix} = 0,$$

Where $\left.\begin{aligned} f_{ij} &= \frac{\partial f_1}{\partial y_l}, \\ g_{ij} &= \frac{\partial g_1}{\partial y_j}, \\ (i,j &= 1, 2) \end{aligned}\right\}$ at $(\alpha^*, \beta^*)$ (4.8)

Have negative real parts. If $\lambda_1, \lambda_2$ are the riits of (4.8) then

$$\lambda_1 + \lambda_2 = (f_{11} - g_{11}) + (f_{22} - g_{22}),$$

$$\lambda_1\lambda_2 = (f_{11} - g_{11})(f_{22} - g_{22}) + f_{21}g_{12}. \quad (4.9)$$

The proof of the following result is an immediate consequence of (4.7)-(4.9)

THEOREM 4.1 *In the CREDIT-DEBIT system (4.1)-(4.6) assumes that the partial derivatives satisfy the following additional conditions:*

$$\frac{\partial g_1}{\partial y_1} > \frac{\partial f_1}{\partial y_1} + \frac{\partial f_2}{\partial y_1},$$

$$\frac{\partial g_2}{\partial y_2} > \frac{\partial f_2}{\partial y_2} + \frac{\partial g_1}{\partial y_2}, \quad \text{Evaluated at } (\alpha^*, \beta^*), \quad (4.10)$$

*Then the steady state $(\alpha^*, \beta^*)$, is (locally) asymptotically stable for the prey- predator system.*

## 11. Delays in reproduction, predation and gestation

A number of authors have proposed and analysed several models of prey- predator systems with time delays. Most of the models have been derived from the so called models of the Kolmogorov-type f (1.2) with time delays incorporated in the average growth rates of the prey and predator species. Volterra [32] proposed the following model for a prey-predator system. The equations are used for the formulation of Credit-debit system of atoms, accounts, economic variable, cosmological variable, astrophysical; variables, and qubit systems for which the model is being applied.

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \delta_1 N_1(t) - \int_{\infty}^{t} F_1(t-s)N_2(s)ds],$$

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 + \delta_2 N_1(t) - \int_{\infty}^{t} F_2(t-s)N_1(s)ds], \quad (5.1)$$

Where $r_1, r_2, \delta_1, \delta_2$ *are positive constant and* $F_1, F_2$ is nonnegative continuous delay kernel defined and integrable on $[0, \infty]$. (5.1) has no predator self-regulating negative feedback. Brelot [6] considered a modified format of (5.1) in the form

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \lambda_1 N_1(t) - \delta_1 N_2(t) - \int_{-\infty}^{t} F_1(t-s)N_2(s)ds], \quad (5.2)$$

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 - \lambda_2 N_2(t) - \delta_2 N_1(t) - \int_{-\infty}^{t} F_2(t-s)N_1(s)ds],$$

Which has term representing intraspecific competition among both predator and prey species. Under suitable condition on the various parameters in (5.2) one can show that (5.2) has a (locally) asymptotically stable steady state. A number of intergrodifferential equation models of the types (5.1) and (5.2) have been investigated by Cushing [12]

Starting from Hutchinson's (1948)[19] delay logistic equation, May[23] has proposed the following system

$$\left.\begin{aligned} \frac{dN_1(t)}{dt} &= rN_1(t)\left[1 - \frac{N_1(t-\tau)}{K}\right] - \alpha N_1(t)N_2(t), \\ \frac{dN_2(t)}{dt} &= -bN_2(t) + \beta N_2(t)N_1(t), \end{aligned}\right\} \quad (5.3)$$

Where r, $\tau$, K, $\alpha, \beta, b$ are positive constants; (5.3) contains a single discrete delay; one can modify (5.3) and incorporate a continuously distributed delay in it so that (5.3) can be written as (May [23])

$$\left.\begin{aligned} \frac{dN_1(t)}{dt} &= rN_1(t)[1 - \int_{-\infty}^{t} F(t-s)N_1(s)ds] - \alpha N_1(t)N_2(t), \\ \frac{dN_2(t)}{dt} &= -bN_2(t) + \beta N_1(t)N_2(t). \end{aligned}\right\} \quad (5.4)$$

The first model of a prey-predator type which departs from the Kolmogorov-type formulation is due to Wangersky and Cunningham [33] who considered the model system

$$\left.\begin{aligned} \frac{dN_1(t)}{dt} &= a_1 N_1(t) - b_1 N_1^2(t) - c_1 N_1(t) - N_2(t), \\ \frac{dN_2(t)}{dt} &= a_2 N_2(t) + c_2 N_1(t-\tau) N_2(t-\tau), \end{aligned}\right\} \quad (5.5)$$

Where $a_1, a_2, b_1, c_1, c_2, \tau$ are positive constants; (5.5) means that a duration of $\tau$ time units elapses when an individual prey is killed and the moment when the corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the

case may be( Note the author deals for predator – population.) Recently Cushing [13] has considered a model of the form

$$\frac{dN_1(t)}{dt} = rN_1(t)\left[1 - \frac{N_1(t)}{k} - \alpha N_2(t)\right],$$

$$\frac{dN_2(t)}{dt} = -d_2^*(t) + bN_1(t)\int_0^\infty \delta(a)N_2(t-a)e^{-d*a}\,da, \quad (5.6)$$

Where $r, k, \alpha, d^*, b$ are positive constants and $\delta$ is related to the age dependent fecundity of the predator species.

With the models (5.1)-(5.6) in the background let us now consider a time delayed model similar to (4.1) in the following form:

$$\frac{dN_1(t)}{dt} = f_1\big(N_1(t-\tau_{11})\big) - g_1\big(N_1(t), N_2(t-\tau_{12})\big),$$

$$\frac{dN_2(t)}{dt} = f_2\big(N_1(t-\tau_{21})\big), N_2(t-\tau_{12}) - g_1(N_2(t)), \quad (5.7)$$

Where $\tau_{ij}$ $(i,j=1,2)$ are nonnegative constants with $\tau = \max \tau_{ij}, i,j = 1,2$; the birth and death rate functions $f_1, f_2, g_1, g_2$ satisfy the same conditions as in (4.1)-(4.6). Along with (5.7) we specify the initial conditions as follows:

$$N_1(s) = \phi_1(s) > 0, \quad s \in [-\tau, 0].$$

$$N_1(s) = \phi_1(s) > 0,$$

We3 will assume that in addition to (4.2) we have for (5.7) that

$$\frac{\partial f_2}{\partial N_1} > 0, \frac{\partial f_2}{\partial N_2} > 0, f_2(N_2, N_2) > 0 \; whenever \; N_1 > 0, \; N_2 > 0. \quad (5.9)$$

It is not difficult to show by the methods previously used for the competition system that whenever $N_1(s) > 0$ and $N_2(s) > 0$ for $s \in [-\tau, 0]$, solution of (5.7) exist for all $t \geq 0$ and remain nonnegative for all $t \geq 0$. also $(\alpha^*, \beta^*)$ of (4.5) is a steady state of (5.7).

We will discuss the form and meaning of the delays in (5.7) in the next section. The effects of the various delays in (5.7) on the (local) asymptotic stability of $(\alpha^*, \beta^*)$ can be derived from the following theorem whose proof is quite similar to that of Theorem 3.1 and hence the details of the proof are omitted.

THEOREM 5.1 *In the time delayed prey-predator system(* corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) *) (5.7) let the conditions (4.2)-(4.6), (5.8) and (5.9) hold for the birth and death rates; furthermore suppose the condition (4.10) hold. Then for all nonnegative delays in (5.7), its nontrivial steady state* $(\alpha^*, \beta^*)$ *is (locally) asymptotically stable.*

## 6. Estimation of decay rates

We will first derive a convention sufficient condition for the asymptotic stability of the trivial solution of a linear system of the form

$$\left.\begin{aligned}\frac{dY_1(t)}{dt} &= -m_{11}\,Y_1(t) + b_{11}Y_1(t-\tau_{11}) - m_{12}Y_2(t-\tau_{12}),\\ \frac{dY_2(t)}{dt} &= -m_{22}\,Y_1(t) + m_{21}Y_1(t-\tau_{21}) - b_{22}Y_2(t-\tau_{22}),\end{aligned}\right\} \quad t>0, \quad (6.1)$$

Where$m_{1J}$ , $b_{i1}$ $(i,j=1,2)$ are positive constants while $\tau_{1J}(i,j=1,2)$ are non negative with $\tau=\max(\tau_{1J})(i,j=1,2$.as is shown in the appendix, the following system dominates (6.1*)*

$$\left.\begin{aligned}\frac{dZ_1(t)}{dt} &= -m_{11}\,Z_1(t) + b_{11}Z_1(t-\tau_{11}) - m_{12}Z_2(t-\tau_{12}),\\ \frac{dZ_2(t)}{dt} &= -m_{22}\,Z_1(t) + m_{21}Z_1(t-\tau_{21}) - b_{22}Z_2(t-\tau_{22}),\end{aligned}\right\} \quad t>0, \quad (6.2)$$

$Z_i(s)=|Y_i(s)|, \qquad i=1,2 \;\; for\; s\le 0 \;\; for\; which\; Y_i(s)\; is\; defined.$

Thus the asymptotic stability of the trivial solution of (6.2) implies that of (6.1). it is also shown in the appendix that the solutions of (6.2) are nonnegative for all t $\geq$ 0. we will use the following vector-matrix notation:

$$\left.\begin{aligned}&Z=(Z_1,Z_2)^{\tau}, \qquad \big||Z|\big| = |Z_1(t)| + |Z_2(t)| = Z_1(t)+Z_2(t)\\ &B=\begin{bmatrix} b_{11} & m_{12}\\ m_{21} & b_{12}\end{bmatrix}, \quad \big||B|\big| = max\begin{matrix} b_{11}+ & m_{21}\\ b_{21}+ & m_{12}\end{matrix},\\ &A=\begin{bmatrix}-m_{11} & 0\\ 0 & -m_{22}\end{bmatrix}.\end{aligned}\right\} \quad (6.3)$$

The logarithmic norm or the measure of any square matrix B, denoted by $\mu(B)$is defined as follows:

$$\mu(B)=\lim_{\theta\to 0+}\frac{/\!/I+\theta B/\!/-1}{\theta}$$

($I$ being an identity matrix of the same size as $B$).The matrix measure compatible with the norms in (6.3) is given by

$\mu(B) = \max_J [b_{jj} + \sum_{\substack{i=1 \\ i \neq J}}^{2} | b_{jJ} | ]$ $B=(b_{1J})$.

We refer to Coppel [11]or Vidyasager [31]for a number of properties of matrix measure We have shown in the appendix that the derivative of $H(Z_1,Z_2)(t)$(seeA.10)satisfies the relation

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq [-m_{11}+b_{11}+m_{21}]Z_1(t)+[-m_{22}+b_{22}+m_{12}]Z_2(t)$, and hence we have

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq \{\mu(A) + /\!/B/\!/\}\{ Z_1(t)+ Z_2(t)\}$

If we assume for the matrix norm $\|B\|$ as in (6.3)that

$$\mu(A) + /\!/B/\!/ < 0 \; or \; -\mu(A) > /\!/B/\!/$$

Then from the results of the appendix (Lemmas 5 and 6) it will follow that $Z_1(t)+Z_2(t) \to 0$ as $t \to \infty$ and thus we have the following:

Lemma 6.1 Assume that the Coefficients of the linear system (6.1)satisfy the condition (6.7). Then the trivial solution of(6.1)is asymptotically stable.

When (6.7)holds an estimate for the decay rate of solutions of (6.1)can be obtained as follows A necessary and sufficient condition for the asymptotic stability of the trivial solution is that all the roots of the characteristic equation

Det $G(\lambda,T_{iJ})=0$

Have negative real parts where

$$G(\lambda,T_{1J})=\begin{bmatrix} \frac{\lambda+m_{11}}{0} & \frac{0}{\lambda+22} \end{bmatrix} - \begin{matrix} b_{11}^{\,e-\lambda T_{11}} \\ \end{matrix} \begin{matrix} -m_{12}^{\,e-\lambda T_{12}} \\ -m_{21}^{\,e-\lambda T_{21}} \end{matrix} \quad \begin{matrix} \\ b_{22}^{\,e-\lambda T_{22}} \end{matrix}$$

$=\lambda I - A - B(\lambda,T_{iJ})$(say),

A positive number $\alpha$(if it exists)is said to be a decay rate of solution of (6.1) if and only if every root of (6.8)satisfies the condition

$Re(\lambda)+\alpha < 0$.

When (6.11)holds the solution of (6.1)will be of the order of $e^{-at}$ and hence the supremum of all such $\alpha > 0$ $a$ satisfying (6.11)will be a desirable parameter to estimate However such a supremum of all such (6.11)may not be possible to determine since it has to be done directly from (6.8)-(6,10)(see also Section 3).We will obtain an estimate of $\alpha$ (not necessarily the supremum)I (6.11)by the following procedure.

Replace λ byz-$\mathfrak{a}$ in (6.8)and define

$G(z-\mathfrak{a},T_{IJ})=M(z)=(z-a)I-A-B(z-\mathfrak{a},T_{IJ})$.

By Lemma 6.1 a Sufficient condition for all roots of M(z)=0 to have negative real parts equivalently to satisfy

$\mathrm{Re}(\lambda)< -a$

is the following :

$-\mu(aI+A) > \|B_1(a,T_{ij})\|$,

Where

$$B_1(a,T_{ij})= \begin{pmatrix} b_{11}^{\,e-\lambda T_{11}} & -m_{12}^{\,e-\lambda T_{12}} \\ -m_{21}^{\,e-\lambda T_{21}} & b_{22}^{\,e-\lambda T_{22}} \end{pmatrix}$$

Since $\mu(\mathfrak{a}\, I+A)=\mathfrak{a}+\boldsymbol{\mu}(\boldsymbol{A})$ we have from (6.13) that

$\mathfrak{a}+\mu(A)+\|B_I(\mathfrak{a},T_{IJ})\|< 0$

Is a sufficient condition for Re $(\lambda) < -a$If we denote the left side of (6.15)by $\mathfrak{a}f(\mathfrak{a},T_{Ij})$ we note that

$f(0,T_{Ij})=\|B\|+\mu(A)<0 \;\; if\;(6.7) holds;$

Al so $\partial f/\partial \mathfrak{a} > 0, f$ is continuous and monotonic in $(\mathfrak{a})$;$\to \infty as\ \mathfrak{a}\to\infty$; thus by the continuity of $f$ on $\mathfrak{a}\geq 0$, there exists an interval of the from $[0,^{*})$such that

$f(0,T_{Ij})< 0 \;\; for\ \mathfrak{a}\in[0,\mathfrak{a}^{*})$

And hence

$f(\mathfrak{a},T_{Ij})=0$

Thus an upper bound $\mathfrak{a}^{*}$ of $\mathfrak{a}$ for which we have $f(\mathfrak{a},T_{Ij})< 0\ \ \mathfrak{a}\in[0,\mathfrak{a}^{*})$ can be obtained as the unique (by the monotonicity of $f$ on $\mathfrak{a}$) positive root of$\max \begin{Bmatrix} b_{11}^{\,e-\lambda T_{11}} & -m_{12}^{\,e-\lambda T_{21}} \\ b_{22}^{\,e-\lambda T_{22}} & -m_{21}^{\,e-\lambda T_{12}} \end{Bmatrix} + \mathfrak{a}\text{-}\min \begin{Bmatrix} m_{11} \\ m_{22} \end{Bmatrix} =0$

Where the maximum in (6.18) is taken over$b_{11}, b_{22}, m_{12}, m_{21}$ and $T_{ij}\,(i,=j\ 1,2)$. A numerical procedure can be devised to solve equations of the form (6.18) in concrete cases. If one is interested in a smaller bound on the decay rate it can be obtained as the unique root of

$\|B\| e^{ar} + a+\mu(A)= 0, \tau = \max_{i,J} T_{ij}$ . (6.19)

It will follow from (6.19) than an increase in T

Will reduce the value of the root of (6.19). We remark that since our estimate $a^*$ of dacay rate is based on a sufficient condition of the form (6.7), we can conclude that the solotions of (6.1) will converge at least as fast as $e^{-a*t}$ where $a^*$ is the unique root of either of (6.18) or (6.19).

**Implications**

Authors have formulated a class of two systems model systems staring from the elementary principal that the rate of change of a population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) equals the creation rate minus the destruction rate. This is a departure from the conventional model building in population ecology where one usually starts from the Malthusian formulation, now known as the Kolmogorov formulation in the nonlinear context. In our analysis of the models proposed we have obtained sufficient conditions for the global existence of nonnegative solutions of the model systems (this is trivial in the Kolmogorov formulation) and the other desirable properties such as the existence of suitable steady states. The two classes of models proposed for competition and debt contain intraspecific self – regulation biologists that each systems is under the action of a self-regulating negative feedback force, opinions and arguments differ regarding the mechanism by which such an effect is realized; we refer for more details on this aspect to a collection of articles compiled by Tamarin [30].

Gopalasamy's main result is the following : if a self –regulating negative feedback mechanism (intraspecific oppugn ant) can act without time delay and if such an effect can keep an ecosystem ( a community of one or more species of populations) asymptotically stable ( at least locally about a non trivial steady state, then time delays of any magnitude in any other interaction involving positive or negative feedbacks cannot destabilize such a stable ecosystem ( in a mathematical sense); that is such delays cannot cause a hopf-type bifurcation leading to persistent and undamped fluctuations. The delay-induced bifurcations to oscillations in models have been investigated by many authors. (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Qualitatively result is not new to the literature the effects of delays on otherwise stable and unstable systems have been considered by Cushing[12]where (pages 35-39)he has shown that an otherwise unstable system can be made less unstable by delays, Such an effect of delays on unstable systems appear to have been first proposed by Beddinton and May [2]who have also shows that delays can make stable systems less stable leading to a conclusion "time delays are not necessarily destabilizing" It is also shown by Cushing([12], pages 75-103)that while time delays cannot stabilize(in the mathematical sense)an otherwise unstable two species system in competition or mutualism time delays in there species systems can stabilize otherwise unstable systems (even in a mathematical sense)Cushing and Saleem [14]have formulated a model where delays in the positive feedback can have a stabilizing effect: Hale[18]discusses briefly the stabilizing nature of delays and given an example to show that an increase in delays can make an unstable system stable Thus for a better understanding of the effects of delays on otherwise stable systems one has to estimate the rate of decay of perturbations (or the characteristic return time )in terms of the parameters of the system especially when the delays do not cause a mathematical instability This aspect has been investigated by a procedure slightly different from that of the characteristic equation we have proposed (section6)

method for the estimation of decay rate of perturbations when there are multiple delays Since many species systems with multiple delays cannot be analysed (due to complexity)by the characteristic equation our methods using comparison lead to easily verifiable sufficient conditions The sufficient conditions obtained are in the form of a diagonal dominance type; these conditions can be weakened to quasi-diagonal dominance type as in Ladde [20], Siljak[20]Anderson [1], Lewis and Anderson [21,22](Ladde's result [20] is delay dependent due to the fact that delay appears in self –regulation also).Verification of quasi –diagonal dominance conditions in arbitrary matrices (except for Metzler matrices )is an impossible task and quasi –diagonal dominance is hard to interpret.

If follows from our discussion that those confederation systems which display delay induced instability leading to permanent and undamped oscillations should have one or more of the following aspects:

(1)time delays are in intraspecific interaction terms (and possibly in other terms also): (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

(2) Intraspecific self-regulation terms are absent in density dependent from in one or more populations of the system;

(3) Intraspecific negative feedback effect are not strong enough to maintain the stability of the relevant steady state of the nondelayed system.

We can now infer how the delay induced stability leads to fluctuations; a primary mechanism of such fluctuations can be described at least theoretically as follows : in its attraction of a system to a steady state, intraspecific self-regulation fails to be effective and the system overhoots the steady state and at such higher densities intraspecific competition becomes operationally effective leading the system to understood the steady state and this process is repeated again and again resulting in undamped oscillations about the steady state. It follows from our results that if intraspecific competition is sufficiently strong, the system cannot become oscillatory even when there are time delays, so long as they do not interfere with the strength of the intraspecific self – regulation.

We are thus led to the question of existence and importance of intraspecific competition in population systems. In the introduction to an anthology of articles by several authors, Tamarin [ 30] begins with the statement "all animal populations have the reproductive capacity to increase without limit, yet none di" and in his concluding remarks writes " this leaves us with only one great concept in the area of confederation do main regulation: that all credit-debit are regulated" More specifically this is applicable to all the confederation of (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

In his discussion on the effect of competition on the densities of populations, Nicholson ([ 25], $page$ [ 147] ) has reached the conclution "Intraspecific competition automatically regulates the severity of its action to the requirements of each denizen , provided the inherent resistance of the environment is sufficiently low to permit the species to exist." In order to look for a mechanism by which intraspecific competition plays its role of decreasing growth rate as confederation density increases, Calhoun [ 7] set up an experiment with Norway rates (Rattus norvegicus) and found that increased density leads to increased hostility which disturbs the psychology and physiology causing decreased reproductive output. Many other mechanisms by which intraspecific compition self-regulate populations have been put forward by a number of authors (see Tamarin [30]). This is true of

(confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

If we accept the view that intraspecific competition is really predominant in population regulation, then we can understand how and why time delays destabilize stable system leading to oscillations; time delays render intraspecific competitive regulation territorially (in case of particles we take the field in which it operates) weak so that self –regulatory forces become temporarily insufficient. If we now agree that all population models should incorporate self-regulating intraspecific competition terms in surfeit strength along with other tropic intraspecific competition terms in sufficient strength along with other atrophic interactions then it is not fair to say " time delays destabilize" and we note from our analysis that it is usually a failure in self- regulation which can cause instability rather than delays by themselves.

We will give one example to illustrate our point. Consider a credit-debit system in which the debits is self –regulating by a logistic load as follows: (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

$$\frac{dx}{dt} = r\left(x - \frac{x^2}{K}\right) - axy = b_I(x) - m_I(x,y) \quad (say),$$

$$\frac{dy}{dt} = \beta xy - \delta y = b_2(x,y) - m_2(y) - \quad (say),$$

Where $K, a, \beta, \delta$ are positive constants such that the steady state $x^* \; y^*$ exists where

$$x^* = \frac{\delta}{\beta}; \quad y^* = \frac{r}{a}\left(-\frac{\delta}{K\beta}\right) > 0.$$

In this example the predator self –regulation $(\partial m_2/\partial y)$ is not sufficiently strong in the sense

$$x^* = \left(\frac{\partial m_2}{\partial y}\right)_{(x^* y^*)} \not> \left(\frac{\partial b_2}{\partial y} + \frac{\partial m_1}{\partial y}\right)_{(x^* y^*)}$$

If in such a model there is a time delay in the credit creation rate, such a delay can lead to instability of $(x^* \; y^*)$ of setting up undamped oscillations. If however we modify the growth rate by the inclusion of a density dependent self-regulation so that credit

$$\frac{dy}{dt} = \beta\, xy - \delta y - \eta y^2$$

For some positive constant η then it is possible, depending on η, for such a modified system to remain stable without oscillations even if there is a time delay in the credit creations rate. We have another completely different justification to have $\eta y^2$ in the credit dynamics: for instance if the debit is removed from this cannot happen. While this example serves to illustrate an aspect of our result, our advocacy of the significance is not based on this example.

Gopalasamy has thus derived conclusion that "time delays need not destabilize" in models with several discrete delays. However we have provided in the appendix the necessary tools (Lemmas 3,4,5) to prove all our results when the relevant delays are all (or some) continuously distributed. Some authors (Caswell [9], May [23], Cushing [12] appear to regard continuously distributed dalays as financially/cosmologically credit wise obtained when he subjected algae to a variable nitrate environment and obtained a better fit for his data with a continuous delay. The questions of whether a discrete or continuous delay is more realistic may well depend on what process is being modeled. For instance an increased supply of some resource at some point in time may trigger reproduction at a particular moment in time later irrespective of the conditions of the resource during the intervening period. If the duration of increased resource supply necessary to elicit such a response is short compared to the intervening period before the reproduction occurs, then a discrete time delay may be appropriate. In a large number of cases of animal. Note fungus algae are prey. Here it shall be an account which loses the energy, amount, or matter for that matter. Same result holds good for (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Populations of particles, atoms, molecules, propositions, systems types chosen the past dietary and nutritional history of an animal over a long period plays an influential role in determining the current bahaviour of the dietary and nutritional history of an animal over a long period plays an influential role in determining the current behavior of the animal and in such cases continuous time delays will be appropriate.

In conclusion the following comments are appropriate. The time delays since the current creation rate is usually due to a contribution of those mature animals living a few units of time earlier depending on the type of animal or organism; the time delays appearing in interspecific interactions can be due to the fact that in order to successfully interact and compete interferentially a developmental maturity may be necessary. In the case of predation, delays in the birth rate of prey can be due to delayed maturation while that in the birth rate of a predator is due to gestation delay with delayed fecundity; the delays in the predation rate are again are again as before due to the maturity needed for a predator to successfully organize and succeed in an attack.

In our analysis we have assumed throughout a temporally uniform environment neglecting dispersal of populations and their age structure. We expect our results to hold under suitable conditions both in temporally periodic environments with the steady state of a constant environment being replaced by a periodic solution (nonstationary steady state) and when the effects of dispersion and age structures are considered. It is worthwhile to modify our purely ecological, cosmological, financial, propositional, problematic systems models so as to provide them with a genetic foundation and examine the effects of time delays since forces of selection act with some time delay. Such a generalization will definitely lead to a better understanding of "how confederations regulate themselves".

## Appendix

First we derive a comparison result.

Lemma 1. Consider the delay differential system

$$\left.\begin{aligned}\frac{dY_1(t)}{dt} &= -m_{11}\ Y_1(t) + b_{11}Y_1(t-\tau_{11}) - m_{12}\ Y_2(t-\tau_{12}),\\ \frac{dY_2(t)}{dt} &= -m_{22}Y_2(t) - m_{21}Y_1(t-\tau_{21}) + b_{22}\ Y_2(t-\tau_{22}),\end{aligned}\right\} (A.1),$$

Where the $m_{ij}$, $b_{ij}$ $(i.j = 1,2)$ are positive while $\tau_{ij}$ $(i.j = 1,2)$ nonnegative with $\tau = \max \tau_{ij}$, $i,j = 1,2$. Then

$$\left.\begin{aligned}\frac{d|Y_1(t)|}{dt} &\le -m_{11}\,|Y_1(t)| + b_{11}|Y_1(t-\tau_{11})| + m_{12}\ Y_2(t-\tau_{12})|,\\ \frac{d|Y_1(t)|}{dt} &\le -m_{22}\,|Y_2(t)| + m_{21}|Y_1(t-\tau_{21})| + b_{22}\ |Y_2(t-\tau_{21})|,\end{aligned}\right\} (A.2)$$

Where

$$\frac{d}{dt}\,|\,Y_1(t)\,| = \left\{\begin{aligned}&\frac{dY_1(t)}{dt} \quad if \quad Y_1(t) \ge 0,\\ &\frac{dY_1(t)}{dt} \quad if \quad Y_1(t) < 0,\end{aligned}\right\} \quad i = 1,2. \qquad (A.3)$$

The proof is by direct verification.

LEMMA 2. Consider now (A.2) and the following system

$$\left.\begin{aligned}\frac{dZ_1(t)}{dt} &= -m_{11}\ Z_1(t) + b_{11}Z_1(t-\tau_{11}) + m_{12}\ Z_2(t-\tau_{12}),\\ \frac{dZ_2(t)}{dt} &= -m_{22}Z_2(t) + m_{21}Z_1(t-\tau_{21}) + b_{22}\ Z_2(t-\tau_{22}),\end{aligned}\right\} \qquad (A.4)$$

Together with the initial conditions

$$Z_1(s) > |Y_1(s)|, \quad s \in [-\tau, 0],\ i = 1,2. \qquad (A.5)$$

Then

$$|Y_1(t)| < Z_1(t) \text{ for } t \ge 0 \quad \text{and} \quad i = 1,2. \qquad (A.6)$$

PROOF. Let $U_1(t) = Z_i(t) - |Y_1(s)|$, $i = 1,2$. $t \ge -\tau$. Then we have from (A.2) and (A.4),

$$\left.\begin{aligned}\frac{dU_1(t)}{dt} &\ge -m_{11}\ U_1(t) + b_{11}U_1(t-\tau_{11}) + m_{12}\ U_2(t-\tau_{12}),\\ \frac{dU_2(t)}{dt} &\ge -m_{22}U_2(t) + m_{21}U_1(t-\tau_{21}) + b_{22}\ U_2(t-\tau_{22}),\end{aligned}\right\} t > 0, \quad (A.7)$$

With

$U_1(s) = Z_i(s) - |Y_i(s)| > 0, \quad i = 1,2, s \in [-\tau, 0]$. (A.8)

Consider (A.7)-(A.8) first on the interval $[\tau, 0]$ : It will follow from (A.7) that $U_1(t) \geq U_1(0) \exp[-m_{11}t]$

$+\int_0^t [b_{11}U_1(s-\tau_{11}) + m_{12}U_2(s-\tau_{12})] \exp[-m_{11}(t-s)ds > 0,]$ (A.9)

For all $t \in [0, \tau^*]$ since $U_i(s) > 0$ on $[-\tau^* 0]$ where $\tau^*$ = positive minimum of $(\tau_{11,}\tau_{12,}\tau_{21,}\tau_{22,})$. Similarly $U_2(t) > 0$ for $t \in [0, \tau^*]$. Now by a similar argument it will follow that $U_1(t) > 0, U_2(t) > 0$ for $t \in [\tau^*, 2\tau^*]$. Repeating the above procedure we obtain the result (A.6).

It follows from Lemmas 1 and 2 that in order to prove theorem 3.1 it is enough to prove that $Z_i(t) \to 0$ as $t \to \infty$ for $i = 1,2$ . which we will do as follows.

PROOF OF THEOREM 3.1 Define a Lyapunov-like function

$H(Z_1,Z_2)(t)=Z_1(t)+Z_2(t)$

$+\int_{1-T}^t b_{11} Z_1(s)ds + \int_{1-T}^t m_{12} Z_2(s)ds$

$+\int_{1-T}^t m_{21} Z_1(s)ds + \int_{1-T}^t b_{22} Z_{22}(s)ds$

Since each term of H in (A.10)is nonnegative for t>0, if we show that $H \to 0$ as $t \to \infty$ it will follow from Lemma 2 that

$0 \leq |Y_{1,}(t)| \leq Z_1(t) \to 0$ as $t \to \infty$,

Implying that the perturbations $Y_{1,}$, $Y_2$, governed by (A.1)approach zero as $t \to \infty$, and this will prove our result Calculate the rate of change of H along the solutions of (A.4)so that

$$\frac{dH}{dt} = \frac{\partial H}{\partial t} + \frac{\partial H}{\partial Z_1}\frac{dZ_1}{dt} + \frac{\partial H}{\partial Z_2}\frac{dZ_2}{dt}$$

$=-[Z_1(t)\{m_{11} - (b_{11} + m_{12})\} + Z_2(t)\{m_{22} - (b_{22}m_{12})\}] < 0$

On using (2.10)since $Z_1(t),Z_2(t)$ are positive for all t>0. Thus $H(Z_1, Z_2)(t) \to 0$ as $t \to \infty$ since H is bounded below by zero: it follows from Lemma 6 below that $Z_1(t) \to 0, Z_2(t) \to 0$ as $t \to \infty$ the proof is complete.

The following set of results are intended to be a guide to prove result similar to the above in the case of continuously distributed delays An analogue of Lemma I is the following.

LEMMA 3 Let $K_{ij}(s)$,$s \in [0,\infty)$, $i, j = 1,2$,denote nonnegative continuous functions integrable on $[0,\infty)$ normalised such that

$\int_0^\infty sk_{ij}(s)ds < \infty$; $\int_0^\infty k_{ij}(s)ds=1$, I,j=1,2.

Consider the intergrodifferential system

$\frac{dY_1(t)}{dt} = -m_{11}Y_1(\mathrm{t}) + b_{11}\int_0^\infty k_{\mathrm{ij}}(\mathrm{s})\, Y_1(\mathrm{t}-\mathrm{s})\mathrm{ds}$

$-m_{12}\int_0^\infty k_{12}(\mathrm{s})\, Y_2(\mathrm{t}-\mathrm{s})\mathrm{ds}$

$\frac{dY_2(t)}{dt} = -m_{22}Y_2(\mathrm{t}) + m_{21}\int_0^\infty k_{21}(\mathrm{s})\, Y_1(\mathrm{t}-\mathrm{s})\mathrm{ds}$

$+b_{22}\int_0^\infty k_{22}(\mathrm{s})\, Y_2(\mathrm{t}-\mathrm{s})\mathrm{ds}$

*Where* $m_{ij}, b_{ij}\ (i,j=1,2)$ *are positive constants Then*

$$\frac{d|Y_1(t)|}{dt} \leq -m_{11}|\,Y_1\,(t)\,| + b_{11}\int_0^\infty k_{11}\,(s)|Y_1(t-s)|ds,$$

$$+m_{12}\int_0^\infty k_{12}\,(s)|\,Y_2(t-s)|ds,$$

$$\frac{d|Y_2(t)|}{dt} \leq -m_{22}|\,Y_2\,(t)\,| + m_{21}\int_0^\infty k_{21}\,(s)|Y_1(t-s)|ds,$$

$$+b_{22}\int_0^\infty k_{22}\,(s)|\,Y_2(t-s)|ds,$$

*Where* $d|Y_1(t)||/dt$ *is defined by* (A.3).

Proof is again by direct verification.

The proof of the following lemma is similar to that of Lemma 2 above.

Lemma 4. *Consider (A.14) and the following intergrodifferential system:*

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1\,(t) + b_{11}\int_0^\infty k_{11}\,(s)Z_1(t-s)ds$$

$$+m_{12}\int_0^\infty k_{12}(s)Z_2\,(t-s)ds,$$

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1\,(t) + b_{11}\int_0^\infty k_{11}\,(s)Z_1(t-s)ds$$

$$+b_{22}\int_0^\infty k_{22}\,(s)Z_2\,(t-s)ds,$$

Together with the initial conditions

$$\infty > Z_i\,(s) > |\,Y_i\,(s)|for\ s \in (-\infty, 0).$$

(A.16)

Then $|Y_i(t)| < Z_i(t)for\ t \geq 0, i = 1,2.$

The following lemma will be useful to prove a result analogous to that of Theorem 3.1 corresponding to a class of continuously distributed time delays.

LEMMA 5. Let $Z_1(t), Z_2\,(t)$ be the solutions of (A.15)- (A.16) where $(Y_1\ Y_2)$ corresponds to an arbitrary but fixed solutions of (A.13). Then as a consequence of Lemma & it will follow that

$$0 \leq |Y_i\,(t)|\ Z_i\,(t) \to \infty \leq as\ t \to \infty, \qquad i = 1,2,$$

(A.17)

Whenever

$\left\{\begin{matrix} M_{11} > b_{11}+m_{21}, \\ M_{22} > b_{22}+m_{12} \end{matrix}\right\}$

PROOF. Define a Lyapunov –like function H where

$H(Z_1,Z_2)(t)=Z_1(t)+Z_2(t)+b_{11}\int_0^\infty k_{11}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$

$+m_{12}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$

$+m_{21}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$

$+b_{22}\int_0^\infty k_{22}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$

Where $(Z_1,Z_2)$is any solution of (A.15)-(A.16) for arbitrary $(Y_1,Y_2)$defined on $(-\infty, 0]$ Computing the rate of change of H in (A.19)along the solution of (A.15)we derive that (on using (A.12))

$\left(\frac{dH(Z_1Z_2(t)}{dt} = [Z_1(t)\{m_{11} - (b_{11} + m_{21})\} + Z_2(t)\{m_{22} - (b_{22} + m_{12})\}]\right.$

$< 0$(using (A.18)) for $Z_1(t) > 0$ $and$ $Z_2(t) > 0$.

By Lemma 6 below the result will follow

Lemma 6. Let H(t)=H$(Z_1, Z_2)(t)$ be as defined in (A.10) Suppose

$$m_{11} > b_{11} + m_{21} \text{ and } m_{22} > b_{22} \mid m_{12}$$

Then

$$Z_1(t) + Z_2(t) \to 0 \text{ast} \to \infty.$$

PROOF. Suppose (A.21) dose not hold; then there exists at least one sequence

$\{t_n\} \to \infty$ as $n \to \infty$ such that

$$[m_{11} - (b_{11} + m_{21})]\, Z_1(t_n) + [m_{22} - (b_{22} + m_{12})]Z_2(t_n) > \varepsilon \qquad (A.22)$$

For some sufficiently small positive number $\varepsilon$ and hence

$$\frac{dH(t_n)}{dt} < -\varepsilon, n = 1,2,3, \ldots.$$

By the continuity and differentiability of H with respect to t,

$H(t_n - \varepsilon) - H(t_n) = \frac{dH(t^*)}{dt}(-\varepsilon), \quad t_n - \varepsilon < t^* < t_n$

If $\varepsilon$ is small we will have

$$= \frac{dH(t^*)}{dt}\varepsilon/2$$

So that

$$H(t_n) \leq H(t_n - \varepsilon) - \varepsilon^2/2$$

$$\leq H(t_{n-1}) - \varepsilon^2/2$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2)$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2):$$

$$\leq H(t_0) - n(\varepsilon^2/2): t_0=0.$$

(A.12) and Lemma 5 will imply that $H(t_0) < \infty$ for any $t_0 > 0$. If n is sufficiently large it will follow that $H(t_n) < 0$ since $H(t_0) < \infty$,which contradicts the nonnegativity of H for all $t \geq 0$. Hence $Z_1(t) + Z_2(t) \to 0$ as $t \to \infty$ implying the asymptotic stability of the trivial solution of (A.1)(since $Z_1$ dominates $y_1$)

We note the result of Lemma 5 and 6 will remain valid if (A.18) is replaced by the condition $\mu(A) + /\!/B/\!/ < 0$

Where A and B are defined as in (6.3).

We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems

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# Model For Innocuously Benign Delay In A Periodic System

**Abstract**

Sufficient conditions are obtained for the existence of a unique linearly denomination stable positive periodic solution of a system model of two competitions in a periodic environment with time lags in interspecific interactions. It is shown that if the self- regulating intraspecific interaction effects are strong enough and act without time delays then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a periodic domination. An example in finance will make things clear. Corporate finance for one big plant generates lot of accounts of the workers whose salary shall be disbursed to the debit of cash credit account. Here Cash credit account is a resource as also the one to which debits are raised, a la predator. Some of the accounts from which the credits come are demand bills purchase, usance bills discount, vostro and nostro accounts in respect of foreign bills and Letters of credit. In economics inventories are used to produce output, output is sold to produce income which might be saved in part and recycled in investment. Cosmologically take an example of a Black hole which gormandizes matter and radiates energy. Notwithstanding information paradox, one thing is clear that matter is being created in to radiation. We show how the model could be applied to axiomatic predications, postulation alcovishness, phenomenological methodologies, propositional statement, and other problems in cosmology, philosophy, and even in Green House effects. We study the sine qua non properties of such a system. In fact these models can be applied to such unsolved problems like Yang mills Theory and mass gap, and other problems. As we apply the model, we shall go on dilating the position status and the applicability of the model. We shall not touch the electrical engineering problems which might rise questions like what is charge of a Bank. We may say when a customer enters a Bank; we can say the Bank is charged. These factors are still relativistic and we shall make an in-depth study as we progress. Gopalsamy's research was published in Australian Journal of Mathematical Society vide: B 25349=365 It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 19. Introduction

At the outset we assume that the system satisfies Haimovici is caution. The purpose of this article is to establish a set of sufficient conditions for the existence of a unique linearly asymptotically stable positive periodic solution for a nonautonomous system of the form

$$\frac{dx(t)}{dt} = x(t)f_1\big(t, x(t), y(t-\tau_1)\big),$$

**(1.1)**

$$\frac{dy(t)}{dt} = y(t)f_2\,\big(t, x(t-\tau_2), y(t)\big).$$

In the context of confederation dynamics a model system of the form (1.1) means that the interspecific interactions occur with time delays $\tau_1$ 0, $\tau_2$ 0 in a periodic environment, if $f_1$ $and$ $f_2$ are assumed to be periodic in $t$ with a common period $\omega$ 0. Note the interactions are inter specific.

Although model systems in periodic domains have been investigated by a large number of authors (for some references see Gopalsamy (4], [5]), periodically varying environments in time delayed model systems of two or more interacting species have not been investigated. The need for Multitype system models in periodic environments and time delayed interactions are evident from the extensive literature of that type relating to single species. While it is known that periodic ecosystems can support more species than uniform environments (Gopalsamy [4]), time delays might destabilize the system leading to a possible loss of one or more species. Furthermore, periodic environments in population dynamics have some relevance to natural selection and evolution; we will consider this aspect briefly at the end.

The principal finding of the following analysis of (1.1) is that, under suitable hypotheses on $f_1$ $and$ $f_2$, the system (1.1) has a linearly asymptotically stable unique positive periodic solution for all nonnegative delays $\tau_1, \tau_2$ (while our numerical experiments suggest that the periodic solution of (1.1) is globally asymptotically stable and is independent of the delay parameters $\tau_1, \tau_2$ it is not known how one can prove such a strong result). Since time delays are usually associated with possible destabilization of otherwise stable systems, we can describe the time delays appearing on (1.1) as "harmless".

It has been recently shown by the author (Gopalsamy [6]) that if the self –regulating negative feedback in each denomination is strong enough and acts without time delays, then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a temporally uniform domain. The result of this paper is thus an extension to a periodic environment of the "harmless" nature of interspecific interaction delays in a temporally constant environment mentioned above.

## 20. Existence and uniqueness of periodic solutions

We will now formulate the problem of existence of periodic solutions of the periodic delay differential system

$$\frac{dx_1(t)}{dt} = x_1(t) f_1\big(t, x_1(t), x_2\,(t-\tau_{12})\big),$$

$t > 0,$

**(2.1)**

$$\frac{dx_2(t)}{dt} = x_2(t) f_2\big(t, x_1(t-\tau_{21}), x_2(t)\big),$$

(Where $\tau_{12}$ , $\tau_{21}$ are nonnegative constants) in a form suitable for the application of fixed point theory of positive operators. We make the following assumptions regarding $f_1$ and $f_2$ in (2.1).

($H_1$) The partial (Frechet) derivatives of $f_1, f_2$ satisfy the following:

$$\frac{\partial f_i}{\partial x_j} < 0, \qquad i,j = 1,2 \text{ For all } t > -\tau \text{ where } \tau = \max(\tau_{12}, \tau_{21}). \qquad \textbf{(2.2)}$$

($H_2$) $f_1(t, x_1, x_2) \equiv f_1(t + \omega, x_1, x_2)$ for all $t > -\tau, x_1 \geq 0, x_2 \geq 0.$ **(2.3)**

($H_3$) There exist positive constants $b_{il}$, $b_i^u$, $a_{il}$, $a_i^u$ ( $l = 1,2$) such that

$$b_{il} - a_{il} x_i < f_i(t, x_1, x_2) < b_i^u - a_i^u x_i,$$

$$b_{il}/a_{il} < b_i^u / a_i^u, \qquad \textbf{(2.4)}$$

For $i$= 1, 2 and $(x_1, x_2) \in R_+ \times R_+$. If $\tau = 0$ then one can show that (2.1)-(2.4) will imply that the nondelayed periodic system (2.1) will have an asymptotically stable strictly positive periodic solution (for details see Gopalsamy [5]); this result will become a special case of what we obtain in the following.

We will use the following notation:

$$B_\omega^\tau = \{h: [-\tau, \infty) \to R, h \text{ is continuous}, h(t + \omega) = h(t), t \geq -\tau\},$$

$$B = B_\omega^\tau \times B_\omega^\tau,$$

$$K_1 = \{h \in B_\omega^\tau | h(t) \geq 0, t \in [-\tau, \infty)\},$$

$$K = K_1 \times K_1,$$

$$||h|| = \left\{\sup_{t \in [0,\omega],} |h(t)|, h \in B_\omega^\tau\right\}$$

$$||x_1, x_2|| = \max\{||x_1||, ||x_2||, (x_1, x_2) \in B\},$$

$$G_1 = \{(x_1, x_2) \in K, ||(x_1, x_2)|| \geq b_l = \min(b_{1l}/a_{1l}, b_{2l}/a_{2l})\}$$

$$G_2 = \left\{(x_1, x_2) \in K, ||(x_1, x_2)|| \geq b^u = \max\left(\frac{b_1^u}{a_1^u}, \frac{b_2^u}{a_2^u}\right)\right\},$$

$\bar{G} =$ closure of $G_2$ *in* $B$ and

$\frac{\bar{G}_2}{G_1} =$ complement of $G_1 in\ \bar{G}_2;\ G_1 \subset G_2$ by **(2.4)**

The cone $K = K_1 \times K_1$ induces a partial order in $B$ and we will write $(x_1, x_2) \leq (y_1, y_2)$ if and only if $x_1 - y_1 \in K_1$. we define an operator $N$ on $B$ as follows

$$N(x_1, x_2)(t) = \begin{cases} N_1(t, x_1, x_2) \\ N_2(t, x_1, x_2) \end{cases} \quad \text{for t> } 0, \qquad \textbf{(2.5)}$$

Where

$$N_1(t, x_1, x_2) = \int_0^{\omega} G_{11}(m, \omega, t-s) g_1\left(m, s, x_1(x) x_2(s-\tau_{12})\right) ds$$

$$+ \int_0^1 G_{12}(m, 0, t-s) g_1\left(m, s, x_1(s), x_2(s-\tau_{12})\right) ds,$$

$$N_2(t, x_1, x_2) = \int_0^{\omega} G_{22}(m, \omega, t-s) g_2\left(m, s, x_1(s-\tau_{21}), x_2(s)\right) ds \qquad \textbf{(2.6)}$$

$$+ \int_0^1 G_{21}(m, 0, t-s) g_2\left(m, s, x_1(s-\tau_{21}), x_2(s)\right) ds,$$

For $t > 0$, with

$$g_1\left(m, s, x_1(s), x_2(s-\tau_{12})\right) = m x_1(s) + x_1(s) f_1\left(s, x_1(s), x_2(s-\tau_{12})\right),$$

**(2.7)**

$$g_2\left(m, s, x_1(s-\tau_{21}), x_2(s)\right) = m x_2(s) + x_2(s) f_2\left(s, x_1(s-\tau_{21}) x_2, (s)\right),$$

$$G_{ll}(m, \omega, t-s) = e^{-m\omega}(1 - e^{-m\omega})^{-1} e^{-m(t-s)}, \qquad i = 1, 2,$$

**(2.8)**

$$G_{ij}(m, 0, t-s) = e^{-m(t-s)}, i \neq j, i, j = 1, 2,$$

$$N(x_1, x_2)(t) = \begin{cases} N_1(n\omega + t, x_1, x_2) \\ N_2(n\omega + t, x_1, x_2) \end{cases} for\ t \in [-\tau, 0], \qquad \textbf{(2.9)}$$

*m* being a suitably selected positive constant (more will be said below) while *n* is a positive constant such that $n\omega > \tau$.

We note immediately the following from the definition of $N$: $if\ (x_1, x_2) \in \bar{G}/G_1$ then $N(x_1, x_2)(t)$ 0 and if $N(x_1, x_2)$ is periodic $\omega$ for $t$ 0 then $N\ (x_1, x_2)$ is periodic in *t* with period $\omega$ for $t$ −$\tau$. The existence of at least one positive periodic solution of (2.1)-(2.4) is established by means of a few lemmas and a result of Schmitt [9].

Lemma 1. *A necessary and sufficient condition for* $(x_1(t), x_2(t))$ , $t$ −$\tau$ *to be a periodic solution of (2.1)-(2.4) of period* $\omega$ *is that* $(x_1, x_2)$ *be a fixed point of N such that*

$$N(x_1, x_2)(t) = \big(x_1(t), x_2(t)\big) \quad and\ (x_1, x_2)\ \in K\ \cap \left(\frac{\bar{G}}{G_1}\right).$$

PROOF. The details of proof are similar to those of Lemma 2 of Busenberg and Cooke [1]. We first verify the continuity of $N(x_1, x_2)$ assuming that $N(x_1, x_2)(t)$ is periodic in *t* of period $\omega$. The continuity of $N(x_1, x_2)$ for $t > 0$ follows from the definition in (2.5)-(2.8) WHICH ALSO PROVIDES $\lim_{t\to 0+} N(x_1, x_2)(\mathrm{t})$ . It will follow from (2.9) and the assumed periodicity of $N(x_1, x_2)\ that\ N(x_1, x_2)$ is continuous

For $t$ −$\tau$. Now we will verify the periodicity of $N(x_1, x_2)$ in $t$. A simple computation leads to

$$N_1(t + w, x_1, x_2) - N_1(t, x_1, x_2) = e^{-mt} F_1(t),$$

**(2.10)**

Where

$$F_1(t) = \int_0^{\omega} G_{11}\ (m, \omega, t - s\ ) g_1\big(m, s, x_1(s), x_2(s - \tau_{12})\big) ds$$

$$+ \int_0^{\omega} G_{12}(m, 0, t - s) g_1\ \big(m, s, x_1(s), x_2(s - \tau_{12})\big) ds$$

$$+ \int_{\omega}^{\omega + t} G_{12}(m, 0, \omega - s) g_1\ \big(m, s, x_1(s), x_2(s - \tau_{12})\big) ds$$

$$- \int_0^{\omega} G_{11}(m, \omega - s) g_1\ \big(m, s, x_1(s), x_2(s - \tau_{12})\big) ds$$

$$+ \int_0^t G_{12}(m, 0, -s) g_1 \left(m, s, x_1(s), x_2(s - \tau_{12})\right) ds$$

**(2.11)**

Which on simplification using the periodicity of $f_1$ shows that $F_1(t) \equiv 0$ whenever $x_1, x_2, \in B$, thus implying the periodicity of $N_1$ in $t$ of period$\omega$. A similar argument for $N_2$ proves that $(x_1, x_2) \in B \Rightarrow N(x_1, x_2) \in B$.

Now suppose $(x_1, x_2)$ is a fixed point of *N and let* $(x_1, x_2) \in K \cap (\bar{G}_2/G_1)$. Then it is easily seen that $(x_1, x_2)$ is a solution of (2.1) and since $(x_1, x_2) \in K \cap (\bar{G}_2/G_1)$ it will follow that $(x_1, x_2)$ is a periodic solution of (2.1).

Conversely, let $(x_1, x_2)$ be a periodic solution of period $\omega$ of (2.1). Then $(x_1, x_2)$ is necessarily a periodic solution of

$$\frac{dx_1(t)}{dt} + mx_1(t) = g_1\left(m, t, x_1(t), x_2(t - \tau_{12})\right)$$

**(2.12)**

$$\frac{dx_2(t)}{dt} + mx_2(t) = g_2\left(m, t, x_1(t - \tau_{21}), x_2(t)\right),$$

For any arbitrary real constant *m* and hence for any fixed constant $m > 0$. It follows from (2.12) that any periodic solution of (2.1) and hence of (2.12) is a fixed point of $N(x_1, x_2)$. We claim that such a fixed point is in $K \cap (\bar{G}_2/G_1)$. Suppose $(x_1, x_2) \notin K \cap (\bar{G}_2/G_1)$. the following possibilities can arise:

(xix) Either $x_1, x_2$ or both $x_1$ *and* $x_2$ may become negative for some value of *t* and it will follow from (2.1) that subsequent value of $x_1, x_2$ or $x_1$ *and* $x_2$ will be negative and hence from (2.4) and (2.1) their derivatives will be negative for all subsequent *t* , implying that $x_1, x_2$ or $x_1$ *and* $x_2$ cannot be periodic.

(xx) While $(x_1, x_2) \in K \subset B$ there may be a value of *t*, say $t^*$, such that $x_1(t^*) > b_1^u/a_1^u$ or $x_2(t^*) > b_1^u/a_1^u$ and in such a case $x_1(t) or\ x_2(t)$ cannot be periodic since respectively $dx_1(t^*)/dt < 0\ or\ dx_2(t^*)/dt < 0$, which will imply that the values of $x_1(t^*)\ or\ x_2(t^*)$ cannot be repeated from below.

(xxi) Similarly if there is a $t^*$ such that $0 < x_1(t^*) < b_{1l}/a_{1l}$ *or* $0 < x_2(t^*) < b_{2l}/a_{2l}$ then such solutions cannot be periodic since $dx_1(t^*)dt > 0$ or $dx_2(t^*)dt > 0$, which will mean that the values of $x_1(t^*)$ or $x_2(t^*)$ cannot be repeated from above.

Thus every periodic solution of (2.1)-(2.4) is a fixed point of *N* in $K \cap \left(\frac{G_2}{G_1}\right) \subset B$.

**LEMMA 2.** *Let N be as in (2.5)-(2.6). Suppose* $\lambda \in (0,1) and \ \lambda N(x_1, x_2) = (x_1, x_2)$ *for some* $((x_1, x_2) \in K - \{0\}$. *then there exist* $\alpha_1^* > 0, \alpha_2^* > 0$ *such that*

$$0 < x_1(t) \leq \alpha_1^* < b_1^u / a_1^u,$$

$$0 < x_2(t) \leq \alpha_2^* < b_2^u / a_2^u,$$

*for* $t \in [-\tau, \infty)$.

PROOF. If $\lambda N(x_1, x_2) = (x_1, x_2)$ then from the definition of N we get for all $t > 0$ and $(x_1, x_2) \in K - \{0\}$,

$$\frac{dx_1(t)}{dt} = \lambda x_1(t) f_2(t, x_1(t), x_2(t - \tau_{12})) + x_1(t) m(\lambda - 1)$$

**(2.13)**

$$\frac{dx_2(t)}{dt} = \lambda x_2(t) f_2(t, x_1(t - \tau_{21}), x_2(t)) + x_2(t) m(\lambda - 1)$$

For all $t \geq 0$ such that $x_1(t) \geq b_1^u/a_1^u$ *or* $x_2(t) \geq b_2^u/a_2^u$ we have $\frac{dx_1}{dt} < 0$ *or* $\frac{dx_2}{dt} < 0$; but these are impossible since $x_1, x_2$ are periodic in t. Thus by the continuity and periodicity of $x_1, x_2$ we have $x_1(t) \leq \alpha_1^* < b_1^u/a_1^u$, $x_2(t) \leq \alpha_2^* < b_2^u/a_2^u$ for all $t \geq 0$; by the definition and periodicity the same will follow for all $t \geq -\tau$. A similar argument holds when $x_1(t) < 0$ or $x_2(t) < 0$, noting that in such a case $dx_1/dt > 0$ or $dx_2/dt > 0$, again contradicting the periodicity of $(x_1, x_2)$ in *t*.

LEMMA 3. *The operator N maps K∩ $(G_2/G_1)$ in to K and is completely continuous*

PROOF. From the definition of N and (2.4) we derive that

$$N_1(t,x_1,x_2) = \int_0^{\omega} G_{11}\,(m,\omega,t-s)x_1(s)[m+b_{1l}-a_{1l}x_1(s)]ds$$

$$+\int_0^{\omega} G_{12}\,(m,0,t-s)x_1(s)[m+b_{1l}-a_{1l}x_1(s)]ds. \quad \textbf{(2.14)}$$

For $(x_1,x_2) \in K \cap (G_2/G_1)$, we have $0 < x_1 < b_1^u/a_1^u$ and hence we get $N_1(t,x_1,x_2) > 0$ for all $t$ by a proper choice of the positive constant say $m_1$.

We make a similar choice say $m_2$ for *m* for the case $N_1(t,x_1,x_2)$ and then choose *m* = max $(m_1,m_2)$. Thus we have

$$N: K \cap (G_2/G_1) \to K \qquad for\ t,$$

And by the periodicity of N we have the same for $t-\tau$.

To prove the complete continuity of N we proceed as follows. We know that if $(x_1,x_2) \in B$ then $N(x_1,x_2) \in B$ and so we can restrict *t* to the set $[0,n\,\omega], n\omega > r$. We write $N = L \circ Q$ where *L* is the linear operator, $L : B \to B$, $L(u_1,u_2)(t)$

$$= \int_0^{\omega} G_{11}(m,\omega,t-s)\begin{bmatrix} u_1 & (s) \\ u_2 & (s) \end{bmatrix} ds, + \int_0^{t} G_{12}(m,0,t-s)\begin{bmatrix} u_1 & (s) \\ u_2 & (s) \end{bmatrix} ds \quad \textbf{(2.15)}$$

And Q is the nonlinear operator

$$Q(u_1,u_2)(t) = \begin{matrix} g_1(m,t,u_1(t),u(t-\tau_{12})) \\ g_2(m,t,u_1(t-\tau_{21}),u_2(t)) \end{matrix} for\ t>0,$$

**(2.16)**

$$Q(u_1,u_2)(t) = Q\,(u_1,u_2)(n\omega+t)\quad for\ t\in[-\tau,0].$$

It can be found that

$$||L(u_1,u_2)(t)|| \leq max\left\{\frac{m+b_1^u}{m},\frac{m+b_2^u}{m}\right\}||u_1,u_2||,$$

**(2.17)**

$$\left\|\frac{d}{dt}L(u_1,u_2)\right\| \leq max\left\{\frac{m+b_1^u}{m},\frac{m+b_2^u}{m}\right\}||u_1,u_2||.$$

The complete continuity of the linear map $L$ follows from (2.17). Now since Q is continuous and maps bounded sets into bounded sets it will follow that $L \circ Q$ is completely continuous.

To prove the existence of periodic solution of (2.1)-(2.4) we will use the following result of Schmitt [9].

**THEOREM (Schmitt).** *Let $H_1$ and $H_2$ be nonempty bounded open neighborhoods of the origin in a Banach space B with a positive cone K such that $H_1 \subset H_2$ and let N: $K \cap (H_2/H_1) \to K$ be completely continuous and which satisfies the following conditions:*

***(m)*** *There exists $(z_1,z_2) \in K \subset B, \|(z_1,z_2)\| = 1$ and a $\mu$ satisfying $\mu > sup\{\|(x_1,x_2) - N(y_1,y_2)\|, (x_1,x_2) \in K \cap H_1, (y_1,y_2) \in K \cap \partial H_2\}$ such that all solution $(x_1,x_2) \in K \cap (H_2/H_1)$ of $(x_1,x_2) = N\,(x_1,x_2) + \lambda(z_1,z_2) with\ 0<\lambda<\mu\ satisfy\ (x_1,x_2) \notin \partial H_1$.*

*(n) All solution $(x_1,x_2) \in K \cap (H_2/H_1)$ of $(x_1,x_2) = \lambda\, N(x_1,x_2), 0<\lambda<1\ satisfy\ (x_1,x_2) \notin \partial H_2$.*

*Then N has at least one fixed point in $K \cap (H_2/H_1)$.*

We can now present our existence result.

**THEOREM 1.** *Assume that* $f_1, f_2$, *of* (2.1) *satisfy* (2.2)-(2.4). *then there exists at least one positive periodic solution of* (2.1)-(2.4).

PROOF. By Lemma 1, it is enough to show that the operator $N: K \cap (G_2/G_1) \to K \subset B$ has at least one fixed point in $G_2/G_1$ and we will do this by applying the result of Schmitt [9]. First we let

$$H_2 = \{(x_1, x_2) \in B | \, ||x_1, x_2|| < \max(b_1^u/a_1^u, b_2^u/a_2^u)\} \equiv G_2,$$

$$H_1 = \{(x_1, x_2) \in B | \, ||x_1, x_2|| < \min(b_{1l}/a_{1l}, b_{1l}/a_{1l})\} \equiv G_1.$$

By Lemma 3, $N: K \cap (H_2/H_1) \to K$ is completely continuous. Choose $\mu = \infty$ and $(z_1, z_2) \equiv (1.1)$. Then all solution $(x_1, x_2) \in K \cap (H_2/H_1)$ of $(x_1, x_2) = N\,(x_1, x_2) + \lambda(z_1, z_2), 0 < \lambda < \infty$, satisfy (2.1) and hence by their periodicity $(x_1, x_2) \notin \partial H_1$. The condition (a) of Schmitt's theorem is satisfied since $\mu = \infty$ by choice and $(x_1, x_2)$ and $(z_1, z_2)$ are periodic. By Lemma 2 the condition (b) of Schmitt's theorem is satisfied. Thus there exists at least one fixed point of N in $K \cap (H_2/H_1)$ and hence there exists at least one positive periodic solution of (2.1)-(2.4) and the proof is complete.

We remark that so far no restriction on the magnitudes of the delays $\tau_{12}, \tau_{21}$ has been imposed except that they be nonnegative.

## 21. Uniqueness of the periodic solution

Applications of fixed point theorems to prove the existence of solutions of nonlinear problems usually suffer from the drawback that the uniqueness of solutions has to be proved often by other methods. We will use the following result of Smith [11] which is a slight variation of his Theorem 3.5 (see also Busenberg and Cooke [1]).

THEOREM (Smith). *Let N be an operator on a cone K in the Banach space B and let* $K^* = \{(x_1, x_2) \in K \cap B, 0 \; x1,x2 \in C*$ *where* $C*$*some positive constant is. Assume the following:*

(s) If $(x_1, x_2) \in K^* - \{(0,0)\}$ and $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

| | |
|---|---|
| (t) | N is monotonic nondecreasing on $K^*$. |
| (u) | If $(x_1, x_2)$ is in the interior of $K^*$ and if $0 < \lambda < 1$ then there exists an $\eta > 0$ depending on $(x_1, x_2)$ and $\lambda$ such that N $(\lambda x_1, \lambda x_2) \geq \lambda (1 + \eta)$ N$(x_1, x_2)$. Then N has at most one fixed point different from $(0,0)$ in $K^* \subset B$. |

The cone K= $K_1 \cap K_1$ where $K_1 = \{h \in B_\omega^\tau : h(t) \geq 0, t \in [-\tau, \infty)\}$ is slightly too wide in that solutions of the form $\{x_1(t), 0\}$ and $\{0, x_2(t)\}$ for (2.1) are also possible. We choose a narrower cone $K^*$ in $B$ as follows: define

$$S = \{(x_1, x_2) \in B : b_{1l}/a_{1l} \geq \|x_1\| \geq b_1^u/a_1^u,\ b_{2l}/a_{2l} \geq \|x_2\| \geq b_2^u/a_2^u,\}, \quad \textbf{(3.1)}$$

And let

$$K^* = \{\lambda(x_1, x_2) \in B : (x_1, x_2) \in S, \lambda \geq 0\}. \quad \textbf{(3.2)}$$

Since $\hat{G}_2 / G_1 \subset K^*$, it is not difficult to see that N: $\hat{G}_2 / G_1 \to B$ has at least one fixed point which is a periodic solution of (2.1)-(2.4). We will prove that this solution is unique by applying the result of Smith cited above. The following lemmas are needed first.

LEMMA 4. If $(x_1, x_2) \in K^* - \{0.0\}$ and if $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

PROFF. If $(x_1, x_2) = N(x_1, x_2)$ and $(x_1, x_2) \in K^* - \{0,0\}$, then $(x_1, x_2)$ is a periodic solution of (2.1) –(2.4) and hence we have

$$\begin{aligned} (b_{1l} - a_{1l}x_1)x_1 &< \frac{dx_1}{dt} < (b_1^u - a_1^u x_1)x_1, \\ (b_{2l} - a_{2l}x_2)x_2 &< \frac{dx_2}{dt} < (b_2^u - a_2^u x_2)x_2, \end{aligned} \quad (\mathbf{3.3})$$

And hence any nonzero periodic solution $(x_1, x_2)$ of (2.1)-(2.4) has to be such that

$$b_{1l}/a_{1l} < x_1(t) < b_1^u - a_1^u,$$

$$b_{2l}/a_{2l} < x_2(t) < b_2^u - a_2^u,$$

For $t > -\tau$, which proves the lemma.

LEMMA 5. The mapping $N$ defined by (2.5)-(2.9) is monotone nondecreasing on $K^*$.

PROOF. Let $(x_1, x_2) \in K^*$, $(y_1, y_2) \in K^*$ such that $(x_1 - y_1, x_2 - y_2) \in K^*$ and $(x_1, x_2) \neq (y_1, y_2)$. Then $x_1 \geq y_1$, $x_2 \geq y_2$. From the definition of $N$ we derive that

$N_1$ (t, $x_1, x_2$) – $N_1$ (t, $y_1, y_2$)

$$= \int_0^{\omega} G_{11}(m, \omega, t-s)\{m[x_1(s)] + x_1(s) f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s) f_1(s, y_1(s), y_2(s-\tau_{12}))\} ds$$

$$+ \int_0^t G_{12}(m, 0, t-s)$$

$$\times \{m[x_1(S) - y_1(s)] + [x_1(s) f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s) f_1(s, y_1(s), y_2(s-\tau_{12}))]\} ds \quad \textbf{(3.4)}$$

With a similar relation for $N_2(t, x_1, x_2) - N_2(t, y_1, y_2)$. Since $f_1$ is continuous, $f$ is bounded on $K^*$; also it will follow from $(x_1, x_2) \neq (y_1, y_2)$ that

$$0 \not\equiv x_1(s) - y_1(s) \geq 0 \text{ and } 0 \not\equiv x_2(s) - y_2(s) \geq 0;$$

A sufficiently large positive constant $m$ will make the right side of (3.4) positive.A similar argument will lead to $N_2(t, x_1, x_2) - N_2(t, y_1, y_2) > 0$ perhaps with a larger constant $m$. Finally we can choose the bigger of these two constants $m$ to yield

$$N(x_1, x_2) - N(y_1, y_2) \geq 0 \text{ for } (x_1 - y_1, x_2 - y_2) \in K^* - \{(0,0)\}.$$

The monotonicity of $N$ on $K^*$ follows.

LEMMA 6. Let $(x_1, x_2) \in$ interior of $K^*$ and let $\lambda$ be $a$ constant such that $0 < \lambda < 1$. Suppose $f_1, f_2$ satisfy (2.2) - (2.4). Then there exists $a$ Positive constant $\eta$ such that

$$N(\lambda x_1, \lambda x_2) \geq \lambda(1+\eta) N(x_1, x_2). \quad \textbf{(3.5)}$$

PROOF. A direct computation leads to

$$N_1(t,\ \lambda x_1,\ \lambda x_2) - \lambda N_1(t, x_1,\ x_2)$$

$$= \lambda \int_0^{\omega} G_{11}\ (m, \omega, t-s) g_1\big(s, x_1(s),\ x_2(s-\tau_{12})\big) H_1(s, \lambda,\ x_1,\ x_2) ds$$

$$+ \lambda \int_0^{t} G_{11}\ (m, 0, t-s) g_1\big(s, x_1(s),\ x_2(s-\tau_{12})\big) H_1(s, \lambda,\ x_1,\ x_2) ds$$

**(3.6)**

Where

$$H_1(s, \lambda,\ x_1,\ x_2) = \frac{m\ x_1(s) + x_1(s) f_1\big(s, \lambda\ x_1(s), \lambda\ x_2(s-\tau_{12})\big)}{m\ x_1(s) + x_1(s) f_1\big(s, x_1(s), x_2(s-\tau_{12})\big)} - 1. \qquad (\mathbf{3.7})$$

Since $f_1$ is decreasing (see (2.2)) with respect to $x_1$ and $x_2$ there exists a positive constant say $\eta_1$ such that

$H_1(s, \lambda,\ x_1,\ x_2) \geq \eta_1$ for $0 < \lambda < 1, s > 0, (x_1,\ x_2) \in K^*$, **(3.8)**

And hence

$$N_1(t, \lambda x_1,\ \lambda x_2) \geq \lambda(1+\eta_1) N_1(t, x_1,\ x_2).$$

A similar argument will show that there exists an $\eta_2 > 0$ such that

$$N_2(t, \lambda x_1,\ \lambda x_2) \geq \lambda(1+\eta_2) N_2(t, x_1,\ x_2),$$

And the result follows if we choose $\eta = \min(\eta_1, \eta_2)$.

Now the following uniqueness result is a direct consequence of the result of Smith [11] and the Lemmas 4, 5, 6.

THEOREM 2. *Suppose $f_1, f_2$ of (2.1) satisfy (2.2) – (2.4). Then there exists a unique $\omega$ - periodic strictly positive periodic solution of (2.1).*

## 4. Linear stability of the periodic solution

Let $\{x_1^*(t), x_2^*(t), t \geq -\tau\}$ be the strictly positive unique periodic solution of (2.1). Let $\{x_1(t), x_2(t), t \geq -\tau\}$ be any other solution of (2.1). To study the linear stability of $(x_1^*, x_2^*)$ we let

$$x_1(t) \equiv x_1^*(t)\{1 + X_1(t)\}, \quad x_2(t) \equiv x_2^*(t)\{1 + X_2(t)\}, \quad t \geq -\tau, \tag{4.1}$$

And derive that the perturbations $(x_1, x_2)$ are governed by the linear variational system (after neglecting the non linear terms in $x_1, x_2$)

$$\frac{dX_1(t)}{dt} = -a_{11}(t)\, x_1^*(t) X_1(t) - a_{12}(t) x_2^*(t - \tau_{12})\, X_2(t - \tau_{12})$$
$$\frac{dX_2(t)}{dt} = -a_{21}(t) x_1^*(t - \tau_{21}) x_1(t - \tau_{21}) - a_{22}(t) x_2^*(t) x_2(t), \tag{4.2}$$

For $t > 0$, where

$$\left.\begin{aligned} -a_{11}(t) &= \frac{\partial f_1}{\partial x_1} \\ -a_{12}(t) &= \frac{\partial f_1}{\partial x_2} \end{aligned}\right\} \text{Evaluated at } (t, x_1^*(t), x_2^*(t - \tau_{12})), \tag{4.3}$$

$$\left.\begin{aligned} -a_{22}(t) &= \frac{\partial f_2}{\partial x_1} \\ -a_{22}(t) &= \frac{\partial f_2}{\partial x_2} \end{aligned}\right\} \text{Evaluated at } (t, x_1^*(t - \tau_{21}), x_2^*(t)).$$

To establish the linear asymptotic stability of $(x_1^*(t), x_2^*(t))$ it is enough to prove that of the trivial solution of (4.2) - (4.3) We will do this by means of a combination of a comparison differential system and a Lyapunov function. First we define

$$\frac{d|X_1(t)|}{dt} = \begin{cases} \frac{dX_1(t)}{dt} & if\ X_i(t) \geq 0, \\ -\frac{d\backslash X_1(t)}{dt} & if\ X_i(t) < 0, \end{cases} \quad i = 1,2, \tag{4.4}$$

And derive from (4.2)- (4.3) the following comparison system of differential inequalities:

$$\frac{d|X_1(t)|}{dt} \leq -b_{11}x_1^*(t)|X_1(t)| + b_{12}x_2^*(t-\tau_{12})|X_2(t-\tau_{12})|,$$
$$\frac{d|X_2(t)|}{dt} \leq -b_{21}x_1^*(t-\tau_{21})|X_1(t-\tau_{12})| - b_{22}x_2^*(t)|X_2(t)|, \quad \textbf{(4.5)}$$

For $> 0$, where

$$-b_{11} = \max_{t\in[-\tau,0]}\{-a_{11}(t)\},$$

$$b_{12} = |\min_{t\in[-\tau,0]}\{-a_{12}(t)\}|$$
$$b_{21} = |\min_{t\in[-\tau,0]}\{-a_{21}(t)\}| \quad \textbf{(4.6)}$$

$$-b_{22} = \max_{t\in[-\tau,0]}\{-a_{22}(t)\}$$

We can now formulate the result as follows. By the periodicity of $a_{ij}$, the constants $b_{ij}$ are independent of $\tau$ (see also (2.2)).

THEOREM 3. *The unique strictly positive periodic solution* $\{x_1^*(t), x_2^*(t), t \in [-\tau,0]\}$ of $(2.1)-(2.4)$ *is linearly asymptotically stable provided*

$$|-b_{11}| > b_{21} \text{ and } |-b_{22}| > b_{12}, \quad \textbf{(4.7)}$$

*Where the constants* $b_{ij}(i,j = 1,2)$ *are as defined in* (4.6).

PROOF. Consider the scalar valued Lyapunov-like function $v(t) = V(t, X_{1,}X_{2,})$where

$$v(t) = V(t, X_{1,}X_{2,}) = |X_{1,}(t)| + |X_{2,}(t)| + \int_{t-\tau_{12}}^{t} b_{12}x_2^*(u)|X_{2,}(u)|du$$
$$+\int_{t-\tau_{21}}^{t} b_{21}x_1^*(u)|X_{1,}(u)|du. \quad \textbf{(4.8)}$$

Since $x_1^*, x_2^*$ are strictly positive, it is found from (4.7)-(4.) that

$\frac{dv}{dt} \leq -[(|b_{11}| - b_{21})\, x_1^*(t)|X_{1,}(t)| + (|b_{22}| - b_{22})\, x_2^*(t)|X_2(t)|]$ **(4.9)**

For $t > 0$. It follows from (4.7)-(4.9) that $V(t, X_{1,}X_{2,}) \leq |X_{1,}(t)| + X_{2,}(t)|$ and $dv/dt \leq 0$,which will imply the stability of the trivial solution of (4.2)-(4.3). We claim that (4.9) also implies the asymptotic stability of the trivial solution of (4.2)-(4.3)in the sense that

$$\lim_{t-\infty} [\,|X_{1,}(t)| + |X_{2,}(t)|\,] = 0,$$

**(4.10)**

For all bounded initial conditions $X_{1,}(s), X_{2,}(s), s \in \{\mathcal{T}, 0\}. Suppose (4.10) is\ not$ true Then we can find an increasing sequence $\{t_m\} \to \infty$ as $m \to \infty$ such that

$(\,|b_{11}| - b_{21})\, x_1^*(t_m)|X_{1,}(t_m)| + (|b_{22}| - b_{12} x_2^*(t_m)|X_2(t_m)| > \varepsilon,$

M=1, 2, 3…

**(4.11)**

Using the right derivative *dv/dt* of $v$ we derive that

$$v(t_m - \varepsilon) - v(t_m) \leq -\int_{t_m-\varepsilon}^{t_m} \frac{dv(T)}{dT}\, d\tau.$$

**(4.12)**

If $\varepsilon$ is small enough we have from (4.9)(4.11),(4.12)that $dv(\tau)/dr < (-\varepsilon/2)$ and

$v(t_m - \varepsilon) - v(t_m) \geq \varepsilon^2/2$ and hence

$v(t_m) \geq v(t_m - \varepsilon) - (\varepsilon^2/2)$

$\geq v(t_m - 1)$-$(\varepsilon^2/2)$

$\geq v(t_m - 2)$-$(\varepsilon^2/2)$

$v(t_0)$-m ( $\varepsilon^2/2$)

Thus for sufficiently large $m, v(t_m)$becomes negative which contradicts the Nonnegative nature of $v\ for\ t$ 0. this contradiction shows that $|X_{1,}(\mathrm{t})| + |X_2(\mathrm{t})| \to 0$ as $\mathrm{t} \to \infty$ and the proof is complete.

As an example let us consider the case of two species Lotka- Volterra type competition in a periodic environment for for which we have

$$f_1\big(\mathrm{t}, x_1(\mathrm{t}), x_2(\mathrm{t}-T_{12})\big)1 \quad \left\{ \frac{x_1(\mathrm{t}) + ax_2(\mathrm{t}-T_{12})}{K_1(\mathrm{t})} \right\}$$

$$f_2(\mathrm{t}, x_1(\mathrm{t}-\ {}_{21})x_2(\mathrm{t})) = 1 - \left\{ \frac{\beta x_1(\mathrm{t}-T_{21})+x_2(\mathrm{t})}{K_2(\mathrm{t})} \right\}$$

Where $\alpha, \beta$ are positive constants and $k_{1,}k_{2,}$ are strictly positive continuous periodic functions with a common period Equation (4.6)and (4.7) for this special case lead to

$$\left|\mathrm{axm}_{t>T}\{-a_{11}(t)\}\right| = \max_{t>T}(1/K_1(t)) > \min_{t>T}\big(-a_{21}(\mathrm{t})\big)\Big| = \min_{t}, \beta/K_2(\mathrm{t})$$

$$\max_{t T}\{-a_{22}(t)\}\Big| = \max_{t T}(1/K_2(t)) > \min_{t T}\big(-a_{12}(\mathrm{t})\big)\Big| = \min_{t} \ \alpha/K_1(\mathrm{t})$$

The inequalities reduce in the absence of temporal variations in $k_{1,}k_{2,}$ to the inequality 0< $\alpha,\beta$<1, which is well known for the global asymptotic stability of the positive steady state corresponding to the above example in a constant systemic domain the above example shows also that delays in interspecific interactions have no influence on the local asymptotic stability of the unique periodic solution of the periodic ecosystem A number of cases of this type in a constant environment have been established recently by the author (Gopalsamy [6]).A numerical example is illustrated at the end.

**Consequences**

It can be found that all the conclusions on the existence uniqueness and stability of a periodic solution of (2.1)-(2.4) are valid when $OT_{12}=0$ or $T_{12}=0$ or both $T_{12}=0$ and $T_{21}=0$: thus it will follow that delays in interspecific competition do not destabilize an otherwise stable coexistence (oscillatory or nonoscillatory) of two competing types of interacting particles atoms molecules transactions interaction. In forthcoming work we will extend the analysis of the system (2.1)-(2.4) to a system of n-competing species with a hypothesis similar to that of (2.4) It is however an open problem to investigate the effects of delays in intraspecific competition in a periodic milieu

Pianka ([8],pages 139-140) remarks that "In real population rate of increase competitive abilities and carrying capacities do very from individual to individual with population density in space and time Indeed temporal variation in the environment may often allow coexistence by continually altering the competitive abilities of population inhabiting it Moreover time lags are doubt lessly of some importance in real population Finally a heterogeneous environment may allow real competitors to evolve divergent resource utilization patterns and to reduce interspecific competitive inhibition

To understand the mechanism of selection pressure against competitive inhibition we first write the Malthusian parameters $\mu_1, \mu_2$ of fitness for the two species as follows (see MacArthur[7],page 226-230):

$$\mu_1=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_1\left(t,x_1(t),x_2(t-T_{12})\right)$$

$$\mu_2=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_2(t,x_1(t-T_{21}),x_2(t))$$

Since $\left(x_1(t),x_2(t)\right)\rightarrow\left(x_1^*(t),x_2^*(t)\right)$ for large t we can write the parameters of finess in the from

$$\mu_1=f_1\left(t,x_1^*(t),x_2^*(t-T_{12})\right)$$

$$\mu_2=f_2\left(t,x_1^*(t-T_{12}),x_2^*(t)\right)$$

In these relations we can use

$$x_2^*(t - T_{12}), x_2^*(t) - T_{12}\frac{dx_2^*(t - T_{12} + \theta_1 T_{12}}{dt}$$

$$x_1^*(t - T_{12}), x_1^*(t) - T_{21}\frac{dx_1^*(t - T_{21} + \theta_2 T_{21}}{dt}$$

$\theta_1, \theta_2, \in (0,1)$ and rewrite $\mu_1, \mu_2$ in the form

$$\mu_1 = f_1\left(t, x_1^*(T), x_2^*(t) - T_{12}\frac{dx_2^*(t - T_{12} + \theta_1 T_{12}}{dt}\right.$$

$$\mu_2 = f_2\left(t, x_1^*(T), -T_{21}\frac{dx_1^*(t - T_{21} + \theta_2 T_{21}}{dt}, x_2^*(t)\right)$$

$\theta_1, \theta_2, \in (0,1)$

It is seen that the presence of the delays $T_{12}$ and $T_{21}$ provides additional sources of variability in the finess parameters and thereby in the selection pressure against interspecific competition Since a necessary and sufficient condition for natural selection to act is differential (variable) rates of confederation growth we find that time delays coupled with periodicity of an domain may allow real to evolve divergent resource utilization patterns and to reduce

Interspecific competitive inhibition for some remarks relating to the significance of a periodic environment to sustain the coexistence of a large number of confederation types in the absence of delays we refer to a recent article (Gopalsamy)The following example has been numerically solved and the solutions are displayed graphically for two different values of the delay parameter $\mathcal{T}$:

$$\frac{dx_1(t)}{dt} x_1(t)[1 - x_1(t) + 0.4x_2(t - \mathcal{T})/k_1(t)]$$

$$\frac{dx_2(t)}{dt} x_2(t)[1 - \{0.8\, x_1(t - \mathcal{T}) + x_2(t)\}/k_2(t)]$$

Where $k_1(t) = 20 + \sin(\pi t)$ and $k_2(t) = 2.5 + \sin(\pi t)$

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There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. “All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ....” (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To “deconstruct” is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between “speech” and “writing.” This binary logic functions in an illicit way to establish speech as the means of giving “presence” to the world, while writing is deemed derivative and inferior In Derrida’s sense of “grammatology,” however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of “differance” (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons

and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti

thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification.

## GOPALSAMY'S MODEL APPIED TO COSMOLOGICAL PROBLEMS

## GLOBAL ASYMPTOTIC STABILITY IN AN ALMOST PERIODIC LOTKA- VOLTERRA SYSTEM

### Abstract

**Gopalsamy studied** Sufficient conditions are obtained for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of a Lotka- Volterra system with almost periodic coefficients. Research was published Australian. Math Soc. B.27 (1986) 346-360. Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of decay rate of perturbations. Sole intention, main objective and primary concern of the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are counted in the same vein as that of predator prey. Note that the model is applicable to each and every module. Quintessentially each parameter should have been given different name such as superscript 1,2, etc., as that has been done in the foregoing and a formal concatenation done at the end. Due to spatial restrictions we have not been able to do the line of thought and manner of analysis as has been done hereinbefore. This is the most important aspect of the model. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated

model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 13. Introduction

At the outset we assume that the system. This article is concerned with the derivation of a set of "easily verifiable" sufficient conditions for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of the Lotka- Volterra system

$$\frac{dx_1(t)}{dt} = x_i(t)\left\{b_1(t) - \sum_{j=1}^{n} a_{ij}\,(t)x_j(t)\right\};$$

$$t>t_0;\ t_0 \in (-\infty,\infty);\ i = 1,2,\dots,n, \qquad (1.1)$$

where $b_1, a_{ij}\ (i,j = 1,2,\dots.,n)$ are nonnegative almost periodic functions defined for $t \in (-\infty,\infty)$.

In question satisfies Haimovici's equation based on age wise stratification given in the foregoing. Note that Haimovici's equations can be applied to Lotka-Volterra system term by term.

Let R and $R_n$ respectively denote the set of all real numbers and the $n$ - dimensional real Euclidean space; $R_n^+$ will denote the nonnegative cone of $R_n$. Since almost- periodic functions defined on R are bounded, we define the constants $b_i^l, b_l^u, a_{ij}^l, a_{ij}^u, (i,j = 1,2,\dots,n)$ by the following:

$$\inf_{t\in R}\ b_l(t) = b_i^l; \qquad \sup_{t\in R} b_l\,(t) = b_l^u, \qquad i,j = 1,2,\dots,n.$$

$$\inf_{t\in R} a_{lj}\,(t) = a_{lj}^l; \qquad \sup_{t\in R} a_{lj}\,(t) = a_{ij}^u.$$

We shall study the almost- periodic system (1.1) with the following assumption on the coefficients of (1.1);

$$b_l^l > 0; \qquad a_{ll}^l > 0, \qquad (1.2)$$

$$b_l^l > \sum_{\substack{j=1 \\ j\neq i}}^{n} a_{lj}^u\,(b_j^u \,/\, a_{lj}^l). \qquad (1.3)$$

Since the solutions of (1.1) corresponding to nonnegative initial conditions remain nonnegative subsequently, it will follow that

$$\frac{dx_1}{dt} \geq x_l\{b_l^u - a_{ll}^l x_i\}; \qquad t > t_0, i = 1, 2, \dots, n, \qquad (1.4)$$

As a consequence of which we will have

$$0 < x_l(t_0) \leq b_l^u / a_{ll}^l = x_i^u \Rightarrow x_i(t) \leq b_l^u / a_{ll}^l, \qquad i = 1, 2, \dots, n; t \geq t_0. \qquad (1.5)$$

(1.7) And (1.5) together lead to

$$\frac{dx_1}{dt} \geq x_l \left\{ b_l^u \sum_{\substack{j=1 \\ j \neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right) - a_{ll}^l x_i \right\}, \qquad i = 1, 2, \dots, n ; \qquad (1.6)$$

Now (1.2), (1.3) and (1.6) lead to

$$x_i(t_0 \geq \left( b_i^l - \sum_{\substack{J=1 \\ J \neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right) \right) (1/a_{ll}^u) = x_i^l$$

$$\Rightarrow x_l(t) \geq x_i^l \ for\ t \geq t_0,\ i = 1, 2, \dots, n. \qquad (1.7)$$

14. It is our aim to show that the mode together with other models stated in the foregoing together with haimovicis model is applicable to whole lot of consummative debit – credit systems in cosmology, astrophysics, finance, economics, sociology, anatomy computation. We will not however divide the system studied here into 9 components/ modules due to restrictions on space….

From the foregoing preparation we have the following:

LEMMA 1.1. *Assume that* $b_i, a_{ij} (i, j = 1, 2, \dots, n)$ *are scalar nonnegative almost periodic functions defined on* $\boldsymbol{R}$ *such that (1.2) and (1.3) hold. Then the set S defined by*

$$S = \{x = (x_1, x_2, \dots, x_n) \in \boldsymbol{R}_n | x_* \geq x_l \geq x^*; i = 1, 2, \dots, n\}$$

(1.8)

*Where*

$$x_* = \min_{1} x_l^l ; \qquad x_* = \min_{1} x_i^u \qquad (1.9)$$

*is invariant with respect to (1.1)*

**Appellation, cognomen, compellation, designation, denomination of the conducive, infelicitous and unfortutious** responses to environment change may be influenced by changes in available domain and milieu under consideration. An example is the following study: Authors describe the gravitational degrees of freedom of the Schwarzschild black hole by one free variable. Authors introduce an equation which we suggest to be the Schrödinger equation of the Schwarzschild black hole corresponding to this model. We solve the Schrödinger equation explicitly and obtain the mass spectrum of the black hole as such as it can be observed by an observer very far away and at rest relative to the black hole. Our equation implies that there is no singularity inside the Schwarzschild black hole, and that the black hole has a certain ground state in which its mass is non-zero. **Schrodinger equation of the Schwarzschild black hole Jarmo Mäkelä (Cambridge U E-Print: gr-qc/9602008 | PDF** Particles, axions, atoms, molecules, quarks, Majorana fermions, bosons et al in its confederation capacity and aggregation processes, type interactions and interactions between landscape dynamics. Following bears ample testimony, infallible observatory to the fact there might be changes in the quantum behaviour due to merger of black holes. Following example bears enough testimony, infallible observatory and impeccable demonstration and apodictic evidence as how some of the models could be applied to cosmological problems. This is attributable and ascribable to the fact that the classification criterion is based on the age.

**How gravity kills (e) Schrödinger's cat Theorists argue that warped space-time prevents (e) quantum superpositions of large-scale objects. Elizabeth Gibney 17 June 2015 Propositions are written in model able form and can easy be scheduled.**

(133) If the cat in Erwin Schrödinger's famous thought-experiment behaved according to quantum theory, it would be able to exist in multiple states at once: both dead and alive.

(134) Physicists' common explanation for why we don't (e) see such quantum superpositions — in cats or any other aspect of the everyday world — is interference from (e) the environment.

(135) As soon as a quantum object interacts with (e&eb) a stray particle or a passing field, it picks (e) just one state, collapsing (eb) into our classical, everyday view.

(136) But even if physicists could completely isolate (eb) a large object in (eb) a quantum superposition, according to researchers at the University of Vienna, it would still collapse into (e&eb) one state — on Earth's surface, at least.

(137) "Somewhere in interstellar space it could be that the cat has (e) a chance to preserve quantum coherence, but on Earth, or near any planet, there's (=) little hope of that," says Igor Pikovski. The reason, he asserts, is gravity.

(138) Pikovski and his colleagues' idea, laid out in a paper published in Nature Physics on 15 June1, is at present only a mathematical argument. But experimenters hope to test whether gravity really does collapse (eb) quantum superpositions, says Hendrik Ulbricht, an experimental physicist at the University of Southampton, UK. "This is a cool, new idea, and I'm up for trying to see it in experiments," he says. Assembling the technology to do so, however, may take as long as a decade, he says.

(139) Cinema-goers who saw the film Interstellar are already familiar with the basic principle behind the Vienna team's work. Einstein's theory of general relativity states that an extremely massive object causes clocks near it to (e) run more slowly because (e) its strong gravitational field stretches (e&eb) the fabric of

space-time (which is why a character in the film aged only an hour near a black hole, while (e) seven years passed on Earth). On a subtler scale, a molecule placed nearer the Earth's surface experiences (e&eb) a slightly slower clock than one placed slightly further away.

(140) Because of gravity's effect on (e&eb) space-time, Pikovski's team realised that variance in a molecule's position will also influence (e&eb) its internal energy — the vibrations of particles within (eb) the molecule, which evolve (eb) over time.

(141) If a molecule were put in a quantum superposition of two places, the correlation between position and (e&eb) internal energy would soon cause (eb) the duality to 'decohere' to (e) the molecule taking just one path, they suggest.

(142) "In most situations decoherence is due to (e) something external; here it's as though the internal jiggling is interacting with (e&eb) the motion of the molecule itself," adds Pikovski.

**A practical limit**

(143) No one has yet seen this effect because other sources of decoherence — such as magnetic fields, thermal radiation and vibrations — are (=) typically much stronger and cause (eb) quantum systems to collapse long before gravity becomes an issue. But experimenters are keen to try.

(144) Markus Arndt, an experimental physicist also at the University of Vienna, already tests whether quantum superpositions can be observed for (e) large, though not cat-sized, objects (see Fattening up Schrödinger's cat).

(145) He sends large molecules through (e&eb) a matter-wave interferometer, a system that gives (eb) each molecule a choice of (e) two different paths.

(146) In the classical view, a molecule travels down just one path; a quantum molecule effectively passes through (e&eb) both routes at once and interferes with (e&eb) itself to create (eb) a characteristic wave-like pattern (see image, below).

(147) A quantum interference pattern built up from complex molecules known as phthalocyanines.(deleted due to spatial constraints)

(148) A similar set up could be used to (e) test for gravity's ability to destroy (e) quantum behaviour: by comparing a vertical interferometer, in which the superposition should soon decohere due to (e) the stretching of time in one path versus (e&eb) the other, with a horizontal set up, where (e) the superposition could remain.

(149) Arndt, who has tested the effect for (e) molecules of up to 810 atoms2, points out those large molecules, would be (=) good for testing the gravitational effect because (e) they contain (e) lots of particles contributing to (e) internal energy.

(150) **But not only would the researchers have to further suppress the external environment to reduce (e) other decoherence effects, they would also need to (e) increase the separation of the two paths from micrometres to metres, or else use molecules perhaps 1 million times more massive. "It's certainly very challenging," says Arndt**.

(151) If gravity's effect does limit (e) quantum behaviour on Earth, tests of quantum reality for (e) large objects may eventually have to (e) move into space, says Angelo Bassi, a physicist at the University of Trieste in Italy.

(152) "But from a deep, fundamental point of view, this is nothing new," he says. A gravitational field is merely another environment to interact with, so invoking it does not (e) explain whether quantum behaviour might lead to (eb) classical reality if gravity's influence were mitigated — for example, by doing the experiment in (eb) gravity-free space.

(153) Nor does the effect described by Pikovski and colleagues say anything about quantum gravity: a theory that would unite gravity and (e&eb) quantum mechanics in (eb) a single description, which many researchers are working towards.

(154) "It's an interesting effect, but it's still quantum physics applied to (e&eb) classical general relativity. In that way it doesn't change our picture of the world," adds Bassi. Nature doi:10.1038/nature.2015.17773

Integrated factors these influences because they deal with ambience, mise en scene change or confederation dynamics, but rarely both. Time series of (e) space time suitability models with spatially explicit stochastic models to explore factors that influence the viability of type viability and confederation feasibility under stable and changing scenarios. Methodology could be applied to interrogation propositions and problems which might subsume one with the other notwithstanding that there is a space of singularities and there is space for completion of present from the past and future. An example is Time image movement from Deleuze. We clearly face a stiff challenge in our attempt to proffer Deleuzian treatments of photography as our critical project brushes up against the fact that Deleuze does not invest photography with (e&eb)the creative and re-creative potential that he associates with (e&eb) both cinema and modern painting. In short, while he highly values other visual artistic forms, he seemingly presents photographic texts as stagnate documents or tools that produce certainty, organize bodies and desires, and iterate hackneyed ideas. He even uses photography as something of a foil to demonstrate the innovation of the cinema and the originality of modern painters. In Francis Bacon: The Logic of Sensation (1981), he notes how "photography has taken over the illustrative and documentary role, so that modern painting no longer needs to fulfill this function, which still burdened earlier painters" (10). He later explains that "photographs are ways of seeing, and as such, they are illustrative and narrative reproductions or representations ... . They are what are seen, until finally one sees nothing else." He treats the photo as an instrument for reproducing representations of reality—a device that iterates images until they are ossified as established stories, icons, or even stagnant perceptions. Indeed, like Bergson before him, Deleuze compares photography to what he identifies as a reductionist model of perception: "what we see, what we perceive, are photographs" (74).[4] The photo is, in effect, an always-already passé sensational experience that numbs our sensitivity to the ongoing vitality of life and provides us with an efficient understanding of the dynamism of our world. But it is compelling that Deleuze theorizes both the quotidian nature and the liberatory potential of photography; he treats it as a common, everyday visual experience that, despite its aesthetic and technological limitations, emancipates the modern painter to explore new artistic opportunities. This special issue of Rhizomes repeatedly provides artistic examples and critical strategies that invite us to see photography as replete with the power to unburden and emancipate that Deleuze associates with both film and painting. Indeed, one way to think of this project of reading Deleuze against Deleuze might be to ask: if photography can free painting and painters to pursue novel creative opportunities, can it likewise free itself, photographic practitioners, and even viewers to create and re-create anew? **Rhizomes » Issue 23 (2012) » Michael Kramp Unburdening Life, or the Deleuzian Potential of Photography Michael Kramp Lehigh University** dmk209@lehigh.edu

**Existence of an almost- periodic solution**

DEFINITION 2.1. A function $f(t,x)$,where $f$ is an $m$- vector, $t$ is a real scalar and $x$ is an $n$- vector, is said to be almost periodic in $t$ uniformly with respect to $x \in X \subset R_n$, $if\ f(t,x)$ is continuous in $t \in R\ and\ x \in X$, and if for any $\varepsilon > 0$, it is possible to find a constant $l(\varepsilon) > 0$ such that in any interval of length $l(\varepsilon)$ there exists a $\tau$ such that the inequality

$$||f(t+\tau,x)-f(t,x)|| = \sum_{l=1}^{m} |f_l(t+\tau,x) = f_l(t,x)|$$

$$< \varepsilon \qquad (2.1)$$

is satisfied for all $t \in (-\infty,\infty), x \in X$. The number $\tau$ is called an $\varepsilon-$ translation number of $f(t,x)$.

Let $\{\lambda_j\}$ denote the set of all real numbers such that

$$\lim_{T\to\infty} T^{-1}\int_0^T f(t,x)\exp(-i\lambda_j t)dt \not\equiv 0 \quad for\ x \in X; i = \sqrt{-1}. \qquad (2.2)$$

It is known that when the set X is separable, the set of numbers $\{\lambda_j\}$ in (2.2) is countable. The set $\{\sum_1^N n_j \lambda_j\}$ for all integers N and integers $n_j$ is called the module of $f(t,x)$.

The following result concerned with module containment will be used in the proof of our existence theorem below.

**MODULE CONTAINMENT THEOREM** (Yoshizawa [7], page 18). *Let $f(t,x)$ and $g(t,x)$ be almost periodic in t uniformly for $x \in D \subset \boldsymbol{R}_n$. If for any compact set $S \subset D$ and for any sequence of real numbers $\{\tau_k\}$ having a limit as $k \to \infty$ (the limit being finite or $\pm\infty$) for which $\{f(t+\tau_k,x)\}$ is uniformly convergent on $\boldsymbol{R}\times S, \{g(t+,x)\}$ is also uniformly convergent on $\boldsymbol{R}\times S$ then the module of $g(t+,x)$ is contained in the module of $f(t+,x)$.*

We can now formulate our existence theorem.

**THEORM 2.1** *Assume that the almost-periodic parameters $b_1, a_{1J}$ $(i,j = 1,2,\dots,n)$ of (1.1) satisfy (1.2)-(1.3). Furthermore assume that there exists a positive number $\mu$ such that*

$$\inf_{1\in\boldsymbol{R}} a_{11}(t) > \sum_{\substack{J=1\\J\neq 1}}^{n} (\sup_{t\in\boldsymbol{R}} a_{Jl}(t)) + \mu, \quad i$$

$$= 1,2,\dots,n. \qquad (2.3)$$

*Then the system (1.1) has a strictly positive (component wise) almost-periodic solution, say $w(t) = w_1(t),\dots,w_n(t), t \in \boldsymbol{R}$, whose module is contained in that of*

$$f(t,x) = \{f_1(t,x_1,\dots,x_n), f_2(t,x_1,\dots,x_n),\dots,f_n(t,x_1,\dots,x_n)\}$$

*Where*

$$f_l(t, x_1, \dots, x_n) = x_l \left\{ b_l(t) - \sum_{j=1}^{n} a_{lj}(t) x_j \right\},$$

$$i = 1, 2, \dots n; (x_1, x_2, \dots, x_n) \in S.$$

(2.4)

PROOF. The set $S$ defined by (1.8)-(1.9) is compact in $\boldsymbol{R}_n$. Let $\{\tau_h\}$ be any sequence of real numbers such that $\tau_k \to \infty$ as $k \to \infty$ and

$$\sum_{l=1}^{n} |f_i(t + \tau_k, x_1, \dots, x_n) - f_l(t, x_1, \dots, x_n)| \to 0 \quad \text{as } k \to \infty \qquad (2.5)$$

Uniformly in $(x_1, \dots, x_n) \in S$ for $t \in (-\infty, \infty)$. for any real number $\beta$let $k_0(\beta)$ be the smallest value of $k$ such that $\tau k_0 + \beta \geq t_0$. Since the invariance of $S$ implies that any solution of (1.1)

$$\{x_1(t + \tau_k, t_0, x_1(t_0)), x_2(t + \tau_k, t_0, x_2(t_0)), \dots, x_n(t + \tau_k, t_0, x_n(t_0))\} \in S$$

For all $t \geq \beta$ and $k \geq k_0(\beta)$ whenever $\{x_1(t_0), \dots, n_n(t_0)\} \in S$, our first task in the proof is to show that the sequence

$$\{x_1(t + \tau_k, t_0, x_1(t_0)), \dots, x_n(t + \tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Converges to a continuous bounded function say $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined for $t \in [\beta, \infty)$ such that $w_1(t) \in [x_*, x^*], t \in [\beta, \infty), i = 1, 2, \dots, n$ and the convergence is uniform on all compact subsets of $[\beta, \infty)$. Since $x_*$ and $x^*$ are independence of $\beta$, it is enough to show that

$$\{x_1(t + \tau_k, t_0, x_1(t_0)), \dots, x_n(t + \tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Is a Cauchy sequence on compact subsets of $[\beta, \infty)$.

Let $U$ be any compact subset of $\beta, \infty)$., and let be any arbitrary positive number. Choose an integer $n_0(\varepsilon, \beta) \geq k_0(\beta)$ so large that for $m \geq k \geq n_0(\varepsilon, \beta)$

and all $t \in (-\infty, \infty)$

$$\sum_{i=1}^{n} |b_1(t + \tau_m) - b_1(t + \tau_k)| < (\mu \varepsilon x_*)/(4x^*), \qquad (2.6)$$

$$x^* \sum_{i=1}^{n} \sum_{j=1}^{n} |a_{lj}(t+\tau_m) - a_{lj}(t+\tau_k)| < (\mu\varepsilon x_*)/(4x^*) \quad (2.7)$$

$$(2nx^*/x_*)\exp[-\mu x_*(\beta + \tau_k - t_0)] < (\varepsilon/2x^*) \quad (2.8)$$

Denote $x_i(t, t_0, x_1, (t_0))$ $by$ $x_1(t)$ brevity, and consider a function $v(s)$ defined by

$$v(s) = v\left(x(s), x(s+\tau_m+\tau_k)\right)$$

$$= \sum_{i=1}^{n} |\log x_1,(s) - \log x_1(s+\tau_m+\tau_k)|, \qquad x \in t_0, \quad (2.9)$$

Where $x(t) = \{x_1(t), \dots, x_n(t)\}$ denotes a solution of (1.1) with $\{x_1,(t_0), \dots, x_n(t)\} \in S$. the invariance of S for (1.1), together with the elementary mean value theorem of differential calculus, implies

$$(1/x^*) \sum_{l=1}^{n} |x_l(s) - x_1(s+\tau_m+\tau_k)| \in v(s)$$

$$\in (1/x^*) \sum_{i=1}^{n} |x_l(s) - x_1(s+\tau_m+\tau_k)|. \quad (2.10)$$

Calculating the right derivative $D^+v$ of $v$ along the solutions of (1.1) and simplifying

$$D^+v(s) = \sum_{i=1}^{n} \{sign[x_i(s) - x_i(s+\tau_m-\tau_k)]\}$$

$$\left\{\frac{d[log x_i(s)]}{ds} - \frac{d[log x_i(s+\tau_m-\tau_k)]}{ds}\right\}$$

$$\in \sum_{l=1}^{n} |b_l(s) - b_l(s+\tau_m-\tau_k)|$$

$$+\sum_{l=1}^{n}(-a_{il}(s)|x_i(s)-x_i(s+\tau_m-\tau_k)|+\sum_{\substack{j=1\\ j\neq l}}^{n}a_{ij}(s)|x_j(s)-x_j(s+\tau_m-\tau_k)|)$$

$$+\sum_{i=1}^{n}\sum_{j=1}^{n}|a_{ij}(s)|a_{ij}(s+\tau_m-\tau_k)|\,x_j(s+\tau_m-\tau_k) \qquad (2.11)$$

$$\leq -\mu x_* v\big(x(s),x(s+\tau_m-\tau_k)\big)+\sum_{l=1}^{n}|b_l(s)-b_i(s+\tau_m-\tau_k)|$$

$$+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|. \qquad (2.12)$$

Let us set $s-\tau_k=\sigma$; then we have from (2.6)-(2.7) that

$$\sum_{i=1}^{n}|b_l(s)-b_1(s+\tau_m-\tau_k)|+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|.$$

$$=\sum_{l=1}^{n}|b_l(\sigma+\tau_k)-b_i(\sigma+\tau_m)|$$

$$+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(\sigma+\tau_k)-a_{ij}(\sigma+\tau_m)|$$

$$<(\mu\varepsilon x^*)/(2x^*). \qquad (2.13)$$

Now an integration of (2.12) over $(t_0,t-\tau_k)$ together with (2.13) leads to

$$v\big(x(t+\tau_k),x(t+\tau_m)\big)\leq v\big(x(t_0),x(t_0+\tau_m-\tau_k)\big)\exp[-\mu x_*(t+t_k-t_0)]$$

$$+[(\mu\varepsilon x_*)/(2x^*)]/\mu x_*)$$

$$\leq[(2x_n^*)/x_*]\exp[-\mu x_*(t+t_k-t_0)]+(\varepsilon/2x^*)$$

$< (\varepsilon/x^*)$ On using (2.8), (2.10) (2.14)

For $m \geq n_0(\varepsilon, \beta)$ and for all $t \in \cup \subset [\beta, \infty]$, Thus we have from the definition of $\upsilon$ and (2.14) that

$$\sum_{l=1}^{n} |log x_i\,(t+\tau_k) - log x_l(t+\tau_m)| < (\varepsilon/x^*) \qquad (2.15)$$

And hence using $x_* \leq x_i(t) \leq x^*, i = 1, 2, \dots, n; t \geq t_0$,

$$\sum_{i=1}^{n} |x_i\,(t+\tau_k) - x_i(t+\tau_m)| < \varepsilon \text{ for all } t \in \cup \subset [\beta, \infty], m \geq n_0(\varepsilon, \beta). \qquad (2.16)$$

The existence of a function $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined on $[\beta, \infty]$, such that $w_i(t) \in [x_*, x^*], i = 1, 2, \dots, n; t \in [\beta, \infty]$ follows from (2.16) and the fact that $x_i(t) \in [x_*, x^*], i = 1, 2, \dots, n;$ for $t \geq t_0$. Since $\beta$ is arbitrary, $w$ is in fact

Defined on $(-\infty, \infty)$ and hence we have

$x_i(t + \tau_k) \to w_l(t)$ as $k \to \infty$ uniformly on compact subsets of $(-\infty, \infty)$, $i = 1, 2, \dots, n.$

(2.17)

Our next task is to show that $w: R \to S$ is differentiable and is a solution of (1.1). Since $x(t) = \{x_1(t), \dots, x_n(t)\}$ satisfies (1.1) we have $\sum_{i=1}^{n} |\frac{d}{dt}\, x_i(t+\tau_k) - \frac{d}{dt}\, x_i(t+\tau_m)|$

$$\leq \sum_{i=1}^{n} |f_l\,(t+\tau_k, x_1(t+\tau_k), \dots, x_n\,(t+\tau_k))$$

$$- f_l(t+\tau_m, x_1(t+\tau_k), \dots, x_n(t+\tau_k))| \qquad (2.18)$$

$$+ \sum_{I=1}^{n} |f_l\,(t+\tau_m, x_1(t+\tau_k), \dots, x_n(t+\tau_k))$$

$$- f_l\big(t + \tau_m, x_1(t + \tau_m), \dots, x_n(t + \tau_m)\big)|$$

Where $f_l\ (i = 1, 2, \dots, n)$are as in (2.4). Since $\mathrm{x}(\mathrm{t} + \tau_k) = x_1(t + \tau_k), \dots,\ x_n(t + \tau_k)\ \in \mathrm{S}$ for large $\tau_k$ for $t$ in each compact subset of $(-\infty, \infty)$ there exists an integer $n_1(\varepsilon) > 0$ such that if $m$ ☒ $k$ ☒$n_1(\varepsilon)$, then for any $\varepsilon > 0$,

$$\sum_{i=1}^{n} |f_l\left(t + \tau_k, x_1(t + \tau_k), \dots, x_n(t + \tau_k)\right)$$

$$-f_1\big(t + \tau_m, x_1(t + \tau_k), \dots, x_n(t + \tau_k)\big)| < (\varepsilon/2). \qquad (2.19)$$

Similarly since $x_l(t + \tau_k)\ \in S$ and $x_l(t + \tau_m)\ \in S$ for $i = 1, 2, \dots, n$ there is an integer $n_2(\varepsilon) > 0$ such that for large enough $m, k$, such that $m$ ☒ $k$ ☒$n_2(\varepsilon)$, we have

$$\sum_{l=1}^{n} |f_l\left(t + \tau_m, x_1(t + \tau_k), \dots, x_n(t + \tau_k)\right)$$

$$-f_1\big(t + \tau_m, x(t + \tau_m), \dots, x_n(t + \tau_m)\big)| < (\varepsilon/2) \qquad (2.20)$$

This equation will be very much useful in chemistry if used with the modules mentioned and written in concatenation with each other. It should be borne in mind that the applicability is modulewise and in no case concatenation should be marginalized. In all cases parameters are different which could be represented with different superscripts and subscripts like 1,23,…….etc., **Anal Bioanal Chem. 2004 Sep; 380(1):36-45. Epub 2004 Jul 29. Electronic tongue for pharmaceutical analytics: quantification of tastes and masking effects Legin A et al**

The organoleptic aspects of pharmaceutical formulations affect their acceptability to the patient and hence can have an important effect on concordance with treatment. Objective evaluation of these aspects, particularly the taste of the formulation and the drug substance it contains, is difficult. Whilst volunteer taste panels can be used to good effect their utility is limited, particularly during very early stage development when the toxicological profile of the active pharmaceutical ingredient (API) is yet to be established in detail. A potentiometric "electronic tongue" has

been applied to analyse a variety of 41 individual substances and mixtures of particular interest for pharmaceutical research and development. The electronic tongue (ET) was capable of discriminating between substances with different taste modalities and could also distinguish different substances eliciting the same basic taste; the ET is promising in terms of quantifying the content of each substance and has an ability to detect nuances of the basic taste (e.g. lingering or short-lived). After calibration the electronic tongue was successfully applied to predicting bitterness strength of binary mixtures with a sweetener in terms of "apparent" or "perceived" quinine content. In order to render a formulation palatable it is often necessary to mask the (usually bitter) taste of the API by the addition of masking agents such as sweeteners and flavours. The ET proved capable of distinguishing between formulations with different levels of sweetener and/or flavour in a manner that was consistent with their masking efficiency as perceived by a small human taste panel. A suitably calibrated ET could have the benefit of providing the pharmaceutical formulator with reliable data concerning the taste of the product quickly and with a reduced need to ask volunteers to taste active pharmaceutical samples. Early development activities could be facilitated when human tasting is usually not possible in the absence of the required toxicological data.

Thus if $m \geq n_0(\varepsilon) = max\{n_1, n_2\}$, we have

$$\sum_{l=1}^{n} \left| \frac{d}{dt} x_1(t + \tau_k) - \frac{d}{dt} x_1 ( (t + \tau_m) \right| < \varepsilon \qquad (2.21)$$

Which shows that

$$\lim_{k\to\infty} \frac{d}{dt} x_1(t + \tau_k) \quad exists\ for\ i = 1, 2, \dots, n$$

Uniformly on all compact subsets of $(-\infty, \infty)$. Thus we have

$$\lim_{k\to\infty} \lim_{k\to\infty} \frac{x_i(t + \tau_k + h) - x_i(t + \tau_k)}{h}$$

$$= \lim_{h\to 0} \lim_{k\to 0} \frac{x_l(t + \tau_k + h) - x_l(t + \tau_k)}{h} \qquad (2.22)$$

$$= \lim_{h\to 0} \frac{w_1(t+h) - w_1(t)}{h}, \quad i = 1, 2, \dots, n,$$

Showing that (since the left side of (2.22) exists) $w_i (i = 1, 2, \dots, n)$ is differentiable. Furthermore we have from (2.22) that

$$\frac{dw_l(t)}{dt} = \lim_{k\to\infty} [\, f_i \left(t + \tau_k, x_1(t + \tau_k), \dots, x_n\ (t + \tau_k)\right)$$

$$-f_l\big(t + \tau_k, w_1(t), \dots, w_n(t)\big)$$

$$+ - f_l\big(t + \tau_k, w_1(t), \dots, w_n(t)\big)]$$

$$= -f_l\big(t + \tau_k, w_1(t), \dots, w_n(t)\big); \quad i = 1, 2, \dots, n, \qquad (2.23)$$

By the uniform almost periodicity of $f_1, f_2, \dots, f_n$ on $R \times S$ and hence $w(t) = \big(w_1(t), \dots, w_n(t)\big)$is a solution of (1.1).

In order to complete the proof, we have to show that $w(t) = \big(w_1(t), \dots, w_n(t)\big)$ is an almost- periodic function whose module is contained in that of $\{f_1(t, x_1, \dots, x_n), f_2(t, x_1, \dots, x_n), \dots, f_n(t, x_1, \dots, x_n)\}$ for $(x_1, x_2 \dots, x_n) \in S$. For this purpose it is in fact sufficient (by the module containment theorem) to show that for any sequence of real numbers $\{\tau_k\}$ for which

$$\{f_1(t + \tau_k, x_1, \dots, x_n), f_2(t, + \tau_k, x_1, \dots, x_n), \dots, f_n(t, + \tau_k, x_1, \dots, x_n)\}$$

Converges uniformly on R for all $(x_1, \dots, x_n) \in S$, the sequence $\{w_1\ (t, + \tau_k)\ \dots, w_n(t + \tau_k)\}$ converges uniformly on R where $\tau_k$ tends to a finite number or $\pm\infty$ as $k \to \infty$.

In the case where $\tau_k \to \tau \neq \pm\infty$ as $k \to \infty$, we have $w_l\ (t, + \tau_k) \to w_i(t + \tau), i = 1, 2, \dots, n$ uniformly on R as $k \to \infty$. Let us suppose $\tau_k \to \infty$ as $k \to \infty$. For any $\varepsilon > 0$ there exists an integer $n_0(\varepsilon)$ such that if $m$ ☒$k$ ☒$n_0(\varepsilon)$ then

$$\sum_{l=1}^{n} |b_l\ (t + \tau_k) - b_l(t + \tau_m)| < (\varepsilon/4x^*)x_*, \qquad (2.24)$$

$$x^* \sum_{i=1}^{n}\sum_{j=1}^{n} |a_{ij}(t+\tau_k) - a_{ij}(t+\tau_m)| < \left(\frac{\varepsilon}{4x^*}\right)\mu x_* \qquad (2.25)$$

$$n(2x^*/x_*)\exp[\mu x_* \tau_k] < (\varepsilon/2x^*). \qquad (2.26)$$

Consider the function $v\big(w(s), w(s+\tau_m-\tau_k)\big)$ for $s \in [t, t+\tau_k]$ defined by

$$v\big(w(s), w(s+\tau_m-\tau_k)\big) = \sum_{l=1}^{n} |\log w_l(s) - \log w_l(s+\tau_m-\tau_k)|. \qquad (2.27)$$

Using the fact $w: \boldsymbol{R} \to S$ and $w$ is a solution of (1.1), calculate the right derivative $D^+v \ of \ v$;

$$D^+v\big(w(s), w(s+\tau_m-\tau_k)\big) = \sum_{i=1}^{n}\{\text{sign}[w_l(s) - w_i(s+\tau_m-\tau_k)]\}.$$

$$\left\{\frac{dw_i(s)}{ds}\cdot\frac{1}{w_l(s)} - \frac{dw_l(s+\tau_m-\tau_k)}{ds}\,\frac{1}{w_l(s+\tau_m-\tau_k)}\right\}$$

$$-\mu\sum_{l=1}^{n}|w_l(s) - w_l(s+\tau_m-\tau_k)| + \sum_{l=1}^{n}|b_i(s) - b_l(s+\tau_m-\tau_k)|$$

$$+x^*\sum_{l=1}^{n}\sum_{J=1}^{n}\left|a_{lJ}(s) - a_{lJ}(s+\tau_m-\tau_k)\right| \qquad (2.28)$$

$$-\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+\sum_{i=1}^{n}|b_l(s) - b_l(s+\tau_m-\tau_k)| \qquad (2.29)$$

$$+x^*\sum_{l=1}^{n}\sum_{j=1}^{n}\left|a_{lj}(s) - a_{ij}(s+\tau_m-\tau_k)\right|.$$

If we set $s-\tau_k = \sigma$ then we have from (2.29),

$$D^{+}v\big(w(s),w(s+\tau_m-\tau_k)\big)$$

$$-\mu x_{*}v\big(w(s),w(s+\tau_m-\tau_k)\big)$$

$$+\sum_{l=1}^{n}|b_l(\sigma+\tau_k)-b_l(s+\ \tau_m)|$$

$$+x^{*}\sum_{l=1}^{n}\sum_{j=1}^{n}\left|a_{ij}\,(\sigma+\tau_k)-a_{ij}\,(\sigma+\tau_m)\right| \qquad (2.30)$$

$$-\mu x_{*}v\big(w(s),w(s+\tau_m-\tau_k)\big)$$

$$+[(\varepsilon\mu x_{*})/(4x^{*})]\ +[(\varepsilon\mu x_{*})/(4x^{*})]$$

(2.31)

Integrating (2.31) over $(t,t+T_k)$

$$v(w(t+T_k)w(t+T_m))$$

Another MODEL where the model is directly applicable with the scheduled ones are the case of insurance and corporate finance. Insurance: Mathematics and Economics Volume 42, Issue 1, February 2008, Retrieval of Black–Scholes and generalized Erlang models by perturbed observations at a fixed time Dedicated to Professor André Dubey on the occasion of his 60th birthday Daniel Neuenschwander

S-stable laws on the real line (more generally on Hilbert spaces), associated with some non-linear transformations (so-called “shrinking operations”), were introduced in [Jurek, Z.J., 1977. Limit distributions for truncated random variables. In: Proc. 2nd Vilnius Conference on Probability and Statistics, June 28–July 3, 1977. In: Abstracts of Communications, vol. 3, pp. 95–96; Jurek, Z.J., 1979. Properties of s-stable distribution functions. Bull. Acad. Polon. Sci. Sér. Math. XXVII (1), 135–141; Jurek, Z.J., 1981. Limit distributions for sums of shrunken random variables. Dissertationes Math. vol. CLXXXV]. In [Jurek, Z.J., Neuenschwander, D., 1999. S-stable laws in insurance and finance and generalization to nilpotent Lie groups. J. Theoret. Probab. 12 (4), 1089–1107], the authors interpreted s-stable motions on the real line as limits of total amount of claims processes (up to a deterministic

premium) of a portfolio of excess-of-loss reinsurance contracts and showed that they led to Erlang's model or to Brownian motion. In [Neuenschwander, D., 2000b. On option pricing in models driven by iterated integrals of Brownian motion. In: Mitt. SAV 2000, Heft 1, pp. 35–39], we considered stochastic integrals whose integrand and integrator are both independent Brownian motions, thus modelling a stochastic volatility; as a result we got an analogue of the Black–Scholes formula in this model, confirming a result of Hull and White [Hull, J., White, A., 1987. The pricing of options on assets with stochastic volatility. J. Finance XLII (2), 281–300]. In the present paper, we will look at a common generalization of these processes, namely s-stable motions on the real line perturbed by a stochastic integral whose integrand and integrator are both (not necessarily independent) s-stable motions. The main result will be that if we can observe the distribution of such so-perturbed s-stable motions (together with the values of the perturbing processes) at time t=1, then we can identify the whole model (including the perturbation) among all models with Lévy processes perturbed by an iterated stochastic integral of two Lévy processes (in the Gaussian case) resp. among all models with a compound Poisson process with drift perturbed by an iterated stochastic integral of two compound Poisson processes (in the completely non-Gaussian case if the perturbing processes have no drift) without knowing anything about the history or about its distribution during 0≤t<1. This applies, e.g., to a situation where several assets obey the same model and one can estimate the distribution at time one by looking at the values of all these assets at time t=1. Interestingly enough, it will be convenient to treat the whole matter in the algebraic framework of the so-called Heisenberg group. This is a concept coming in fact from quantum mechanics and is in a certain sense the simplest non-commutative Lie group.

$ e^{-\mu x * T_k}\, v\big(w(t), w(t+T_m-T_k)\big) + (\varepsilon/2x^*)$

$ e^{-\mu x * T_k} \sum_{i=1}^{n} |\log\ w_i(t) - \log\ w_i(t+T_m-T_k)| + (1/\mathrm{x}_*) + (\varepsilon/2x^*)$

n(2x^* / x_*)exp[−μx_* T_k] + (ε/2x^*)

<(ε/x^*).

Since

$$v\big(w(t+T_k)w(t+T_m)\big) \sum_{i=1}^{n} |\ w_1(t+T_k) - w_1(t+T_m))/\ x^*)$$

We have from(2.33)that

$$\sum_{i=1}^{n} | w_1(t+T_k) - w_1(t+T_m)| < \varepsilon \; for \; m \; k \; h_0(\varepsilon)$$

Which show that the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\}$converges uniformly on **R**. By a similar argument the above conclusion is also valid when $T_k \to -\infty \; as \; k \to \infty$. Thus the uniform convergence of the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\} \; as \; k \to \infty \; on$ **R** follows and this completes the proof of the theorem.

**COROLLARY.** Suppose $b_i a_{ij}$ (i, j = 1.2, … , n)in (1.1)are periodic with a common period say $T >$ $0. Then \; if$ (1.2), (1.3)$and$ (2.3)$hold \; the \; system$ (1.1)$will \; have \; a \; periodic \; solution \; w(t) =$ $\{w, (t), \ldots\ldots, w_n(t)\} of \; period \; T \; such \; that$

$$w_1(t+T) = w_1(t), \qquad t \in \boldsymbol{R}: i = 1,2, \ldots, n.$$

**PROOF.** Proof follows from that of Theorem 2.1 if we Choose $T_k = KT(K = 0,1,2, \ldots) and \; note \; that$

$$x_1(t+KT) - w_1(t) \to 0 \; as \; k \to \infty; i = 1,2, \ldots, n; t \in \boldsymbol{R},$$

$$x_i(t+(K+1)t) - w_1(t) \to 0 \; as \; \boldsymbol{K} \to \infty; i = 1,2, \ldots, n; t \in \boldsymbol{R},$$

$$x_i(t+T+KT) - w_1(t+T) \to 0 \; as \; k \to \infty; i = 1,2, \ldots, n; t \in \boldsymbol{R},$$

Implying that

$$w_1(t+T) = w_i(t); \quad i = 1,2, \ldots, n; t \in \boldsymbol{R},$$

## 3. Global asymptotic stability

By definition we say that an almost-periodic solution say $w(t) = \{w_1(t), \ldots, w_n(t)\} of (1.1) is \; globally \; asymptotically$ stable (or attractive) if and only if every other solution x(t)={ $x_1(t), \ldots, x_n(t)$}with $x_1(o) > 0, i = 1.2, \ldots, n$ is defined for all t 0 and satisfies

$$\lim_{t\to\infty} |w_1(t) - x_1(t)| = 0; \quad i = 1,2, \ldots, n.$$

A consequence of such a global asymptotic stability is that there cannot be another strictly positive almost-periodic solution of (1.1)

**THEOREM 3.1** Suppose the almost- Periodic coefficients $b_1$, $a_{ij}$ (i, j = 1,2, … , n)

Satisfy (1.2)and (1.3)and the following

$$\inf_{t\in R} a_{ij}(t) > \left\{ \sum_{\substack{J=1 \\ J\neq 1}}^{n} \sup a_{ij}(t) \right\} + \mu; \qquad i = 1,2, \dots, n$$

For some positive number $\mu$ *Then* (1.1)has an almost-periodic solution (with strictly positive components) which is globally asymptotically stable.

We will omit the proof since the proof is identical with that of Theorem 3.1 of [3]

Under condition (3.2) the almost –periodic solution $w(t) = \{w_1(t), \dots . w_n(t)\}$ of (1.1)enjoys a somewhat strong from of stability even if the coefficients of (1.1)are subjected to a class of perturbations We note (see Yoshizawa [7]page17)that if a function b:R→ $R$ is almost periodic then H(b),the "hull of b" is defined as follows:

H(b)= {g|g:**R**→ $\boldsymbol{R}$ *and* $g(t)\lim_{k\to\infty} b(t + T_k)$ *for* some sequence

$T_k$ of real numbers the limit being uniform in $t \in R$.

By Theorem 2.2 of Yoshizawa [7, page 10], members of the hull of almost periodic functions are themselves almost periodic.

A system of the from

$$\frac{dy_1(t)}{dt} = y_1(t)\left\{B_i(t) - \sum_{J-1}^{n} A_{ij}(t) y_1(t)\right\},$$

*i=1,2,…,n; t*>$t_0 \in R$,

Where

$$B_i \in H(b_i); (A_{ij}), \qquad \text{i, j} = 1,2, \dots, \text{n},$$

is known as a system from the hull of (1.1) As a consequence of the definition of the hull in (3.3)and (3.5)we shall have

$$\inf_{t\in R} B_i(t) = \text{b}_i^l; \quad \sup_{t\in R} B_i(t) = \text{b}_i^l;$$

$$\inf_{t\in R} A_{ij}(t) = \text{b}_i^l; \quad \sup_{t\in R} A_{ij}(t) = \text{a}_{iJ}^u;$$

and hence the coefficients of(3.4)- (3.5)satisfy the same conditions as (1.2),(1.3)and (3.2)It will then follow from theorems(2.1)and (3.1)that solutions of (3.4)-(3.5)with $y_1(t_0) > 0, i = 1,2, \dots, n$ *will have* the following property:

$$y_1(t) > 0 \ \ for\ t \geq t_0; i = 1,2, \dots, n,$$

$$\sum_{i=1}^{n} y_1(t) remains\ bounded\ for\ t \geq t_0$$

**DEFINITION** A Solution $x(t) = \{x_1(t), \dots, x_n(t)\}, t \in t_0 \in \boldsymbol{R}$,of (1.1)is said to be stable under perturbations from the hull of (1.1)if the following holds :let $y(t) = \{(y_1(t), \dots, y_n(t)\}$ be any solution of (3.4)-(3.5)such that $y_1(t_0) > 0$, i=1,2,…,n,and for some positive constants $\beta_1, \beta_2 \beta_1, \leq y_1(t) \leq \beta_2, t \geq t_0, \quad i = 1,2, \dots, n$ Then given any $\varepsilon > 0$,there exists a constant $\delta > 0$ such that

$$\sum_{j=1}^{n} |x_1(t_0) - y_1(t_0)| < \delta$$

$$\sum_{i=1}^{n} \sup_{t\in R} |b_1(1) - B_1(t)| < \delta \qquad = \sum_{i=1}^{n} |x_1(t) - y_i(t)| < \varepsilon$$

$$\sum_{i=1}^{n}\sum_{j=1}^{n} \left| \sup_{t\in R} |a_{ij}(t) - A_{ij}(t)| < \delta \right.$$

For $t \geq t_0$.

We can now formulate the following result which shows that all solution of (1.1)are stable under perturbations from the hull of (1.1)

**THEOREM 3.2**. Assume that the conditions of Theorem (3.1) hold. Then every strictly positive solution of (1.1) ( including its unique almost-periodic solution) is stable under perturbations from the hull of (1.1).

PROOF. Let $x(t) = \{x_1(t), \dots, x_n(t)\}$ and $y(t) = \{y_1(t), \dots, y_n(t)\}$ be any two strictly positive solutions of (1.1) and (3.4) – (3.5) respectively for $t \geq t_0 \in R$. Let $\beta_1, \beta_2$ be constants such that

$$\begin{matrix} \beta_1 \leq x_i(t) \leq \beta_2, \\ \beta_1 \leq y_i(t) \leq \beta_2, \end{matrix} \quad i = 1,2,\dots,n; t \geq t_0 \qquad (3.7)$$

Consider now a function $V(t) = V(x(t)y(t))$ defined by

$$V(x(t)y(t)) = \sum_{i=1}^{n} \left| \log x_1(t) - \log, y_i(t) \right|, \; t \geq t_0 \in R. \qquad (3.8)$$

It is easy to see that

$$\frac{1}{\beta_2} \sum_{i=1}^{n} \left| x_i(t) - y_i(t) \right| \leq V(x(t)y(t))$$

$$\leq \frac{1}{\beta_1} \sum_{i=1}^{n} \left| x_i(t) - y_i(t) \right|, \; t \geq t_0 \in R. \qquad (3.9)$$

Calculating the right derivative $D^+V$ of $V$ and using (1.1) and (3.4) we derive (after simplification) that

$$D^+V\left(x(t),y(t)\right) \leq -\mu\sum_{i=1}^{n}\left|x_i(t)-y_i(t)\right| + \sum_{i=1}^{n}\left|b_i(t)-B_i(t)\right|$$

$$+\beta_2\sum_{j=1}^{n}\sum_{i=1}^{n}\left|a_{ij}(t)-A_{ij}(t)\right|; \quad t \geq t_0,$$

$$\leq -\mu\,\beta_1 V\left(x(t),y(t)\right) + \sum_{i=1}^{n}\left|b_i(t)-B_i(t)\right|$$

$$+\beta_2\sum_{i=1}^{n}\sum_{j=1}^{n}\left|a_{ij}(t)-A_{ij}(t)\right|; \quad t \geq t_0, \tag{3.10}$$

An integration of (3.9) over $[t_0, t]$ with an application of (3.8) leads to

$$\frac{1}{\beta_2}\sum_{i=1}^{n}\left|x_i(t)-y_i(t)\right| \leq V(x(t),y(t)) \leq V(x(t_0),y(t_0))$$

$$+\left[\sum_{i=1}^{n}\sup_{t\in R}\left|b_i(t)-B_i(t)\right| + \beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}\left|a_{ij}(t)-A_{ij}(t)\right|\right]\frac{1}{\mu\,\beta_1},$$

$$t \geq t_0,$$

And hence

$$\sum_{i=1}^{n}\left|x_1(t)-y_1(t)\right| \leq \left(\frac{\beta_2}{\beta_1}\right)\sum_{i=1}^{n}\left|x_1(t_0)-y_1(t_0)\right|$$

$$+\frac{\beta_2}{\mu\beta_2}\left[\sum_{i=1}^{n}\sup_{t\in R}\left|b_i(t)-B_i(t)\right|\right.$$

The construction reason is because is criticized in a number of usage guides: The reason for the long delays was because the costs greatly exceeded the original estimates. One objection to this construction is based on its redundancy: the word because (literally, by cause) contains within it the meaning of reason; thus saying the reason is because is like saying "The cause is by cause," which would never be said. A second objection is based on the claim that because can introduce only adverbial clauses and that reason is requires completion by a noun clause. Critics would substitute that for because in the offending construction: The reason for the long delays in completing the project was that the costs.

$$+\beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}\left|\ a_{ij}(t)-A_{ij}(t)\right|\Big]. \qquad (3.11)$$

Now for any $\varepsilon>0$ if one chooses a $\delta>0$ such that

$$\delta<\min\left\{\frac{\varepsilon}{3}\ \frac{\beta_1}{\beta_2},\frac{\varepsilon}{3}\ \frac{\beta_1}{\beta^2}\mu\right\} \qquad (3.12)$$

Then (3..6) will be satisfied and this completes the proof

## . Jia Li's Presciential Dynamics And Contential Staticity Of Accentuation Attrition Models Applied To Cosmological And Astrophysical And Quantum Mechanical Engineering Problems

**Abstract**: **Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics ET AL ). We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept**

**formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems. Predator prey model literally form bastion, pillar, post, stylobate, sentinel towards the end of dismasting many of the asymptotic and solutional behaviour in finance, economic, and as we shall see in cosmology, quantum computation and other fields.**

## MODEL FORMULATION

Let $p_1(t, a_1)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_1}{dt} = -b_1 Q_1 + g_1(g_1)Q_1$$

Where $b$ is the natural death rate in the absence of credit types and $g(P)$ is the conversion rate function of debits

into credit creations which is assumed to be differentiable and satisfies

$$g_1(0) = 0, \quad g_1(P_1) > 0, \ P_1 \ 1, \qquad \lim_{P_1 \to \infty} g_1(P_1) = (K_1)_1 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_1 + (\rho_a)_1 + D_1(a_1, P_1, Q_1)\rho_1 + 0, \qquad 0 < a_1 \ \infty,$$

$$\rho_1(t,0) = \int_0^\infty n_1(a_1, P_1, Q_1)\rho_1(t_1, a_1)da_1$$

Where $D_1$ is the mortality function and $n_1$ is the fertility function.

Put $_1$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_1}{dt}\right]_1 = x_1\left[\frac{dQ_2}{dt}\right]_1 - x_2\left[\frac{dQ_1}{dt}\right]_1$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_1 - x_4\left[\frac{dQ_2}{dt}\right]_1$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_1 - x_6\left[\frac{dQ_3}{dt}\right]_1$$

Where $x_1, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_1$ and $g_1(P_1)Q_1$

Now when you add the term $(b_1Q_1)$ with $g_1(P_1)(Q_1)$ we are infect see is that $(b_1Q_1)$ is credited incrementally part of $g_1(P_1)(Q_1)$ or instant nearly, when subtracted from $\frac{dQ_1}{dt}$ , it means that $g_1(P_1)(Q_1)$ is slowly eaten by $\frac{dQ_1}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_1 = \mu_1(a_1) + (f_1)_1(P_1)_1 + (f_2)_1(Q)_1 \quad (2)$$

$$n_1 = \beta_1(a_1)h_1(Q_1), \quad (3)$$

Where $\mu_1$the age- dependent debit rate whit is $\mu_1(a_1)$ 0 $\forall a > 0$; $(f_1)_1$ is the density- dependent death rate with $f_1(0) = 0$, $(f_1^1)_1(P_1)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_1)_1(P_1) = (k_1)_1$ ; $f_2$ is the functional response of creditors

to prey with $[f_2]_1(0) = 0, (f_2^1)_1(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_1\ (Q_1) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^1(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_1 + (\rho_a)_1 = -\mu_1(a_1) + (f_1)_1 P_1) + (f_2)_1(Q_1))(\rho_1)_1 \quad (4)$$

$$\rho_1(t_1, 0) = h_1(Q_1)\int_0^\infty \beta_1(a_1)\rho_1(t, a_1)da, \quad (5)$$

$$\frac{dQ_1}{dt} = -b_1 Q_1 + g_1(P_1)Q_1 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_1$ and $Q_1$ of (4-6) satisfy

$$0 \ \rho_1(t, a_1) \ \rho_1^*(t, a_1),$$

$$0 \ Q_1(t) \ Q_1^*(t),$$

Where $\rho_1^*(t, a_1)$ and $Q_1^*(t)$, are solutions of

$$(\rho_t^*)_1 + (\rho_a^*)_1 = -(\mu_1(a_1) + (f_1)_1(P_1^*) + (f_2)_1 Q_1^*))\ \rho_1^*, \quad (7)$$

$$\rho_1^*(t, 0) = \int_0^\infty \beta_1(a_1)\ \rho_1^*(t, a_1)da_1 \quad (8)$$

$$\frac{dQ_1^*}{dt} = -b_1\ Q_1^* + g_1(P_1^*)Q_1^*, \quad (9)$$

With $P_1^*(t) = \int_0^\infty \rho_1^*\ (t, a_1)da_1$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_1^*(t)}{dt} = P_1^*(t)\big(r_1 - (f_1)(P_1^*) - (f_2)_1(Q^*)\big), \quad (10)$$

$$\frac{dQ_1^*(t)}{dt} = Q_1^*(t)(-b_1 + g_1(P_1^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_1(a)e^{-\int_0^a \mu_1(\tau)dt}e^{-ra}\ da_1 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_1^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(t)dt}\ da_1 > 1, \quad (H1)$$

Which guarantees $r_1 > 0$.

*Remark.* If $\beta_1(a) = [\beta_0]_1 a_1 e^{-\alpha_1 a_1}, \mu_1(a_1) \equiv \mu_1$ then the solution of (12) is

$$r_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_1]_1(P_1) = 0, [f_2](Q_1) = \lambda_1 Q_1\ h_1(Q_1) = 1, \quad and\ \ g_1(P_1) = c_1 P_1$$

$$[\rho_1]_1 + [\rho_a]_1 = -(\mu_1 + \lambda_1 Q_1)\rho_1$$

$$\rho_1(t,0) = \int_0^{\infty} [\beta_0]_1 a_1 e^{-\alpha_1 a_1} \rho_1(t,a_1)\, da_1$$

$$\frac{dQ_1}{dt} = -b_1 Q_1 + c_1 P_1 Q_1$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_1}{dt} = \omega_1 P_1 - \lambda_1 P_1 Q_1$$

$$\frac{dQ_1}{dt} = b_1 Q_1 - c_1 P_1 Q_1$$

Where $\omega_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$ THE EQUILIBRIUM $(\rho_0)_1(a_1), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_1(a_1), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_1 + (\rho_a)_1 = -\big(\mu_1(a_1) + (f_1)_1(P_1)\big)\rho,$$

$$\rho_1(t,0) = \int_0^{\infty} \beta_1(a_1)\rho_1(t,a_1)da_1 \qquad (13)$$

The solution of (13) has the form of

$$\rho_1(t,a_1) = Ce_1^{r(t-a_1)} - \int_0^{a} (\tau)d\tau - \int_0^{1} (f_1)_1\big(P_1(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_1(t) = C_1 e^{r_1 t - \int_0^t (f_1)_1(P_1(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_1^{-1}(r)$ if $k_1 > r$.

This is equivalent to

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{(k_1)_1 a_1} da_1 < 1.$$

In fact, if we define

$$R_1(s_1) = \int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau)d\tau} e^{-s_1 a_1}\, da_1$$

$R(s)$ is a decreasing function. $R_1(r_1) = 1$ And $R_1(k_1)_1 < 1$ imply $r_1 < (k_1)_1$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_1(a_1), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{k_1 a_1} da_1 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_1 \equiv \int_0^\infty (\rho_0)_{11}(a_1)da_1$ *satisfies*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{-(f_1)_1(P_0)a_1} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_1(a_1), 0)$

Let $u_1(t, a_1) = \rho(t, a_1) - \rho_0(a_1)$ Then

$$N_1(t) \equiv \int_0^\infty u_1\,(t, a_1)da_1 = \int_0^\infty \rho_1(t, a_1)da_1 - \int_0^\infty (\rho_0)_1(a_1)da = P_1(t) - P_0)_1$$

And

$$(u_1)_1 + (u_a)_1 = (\rho_t)_1 + (\rho_a)_1 - (\rho_0^1)_1(a_1)$$

$$= -\mu_1(a_1) + (f_1)_1(P_1) + (f_2)_1(Q_1))\,(u_1 + (\rho_0)_1) + \mu_1(a_1)(\rho_0)_1 + (f_1)_1(P_0)_1(\rho_0)_1$$

$$= -\mu_1(a_1)u_1 - (f_1)_1(P_1) + (f_2)_1(Q_1))\,u_1 - (\rho_0)_1\,(f_1)_1(P_1) - (f_1)_1(P_0)_1$$

$$+ (f_2)_1(Q_1)), \qquad (15)$$

Where $(\rho_0^1)_1(a_1)$ denotes $d(\rho_0)_1(a_1)/da_1$

Expansion $(f_1)_1(P_1)$ and $(f_2)_1(Q_1)$ at $(P_0)_1$ and 0 yield respectively

$$(f_1)_1(P_1) = (f_1)_1((P_0)_1 + N_1)$$

$$= (f_1)_1(P_0)_1 + (f_1^1)_1(P_0)_1N_1 + o_1(N_1), \qquad (16)$$

$$(f_2)_1(Q_1)$$

$$= (f_2^1)_1(0)Q + o_1(Q_1) \qquad (17)$$

Write $f_1(P_0) = f_{10}, f_1^{'}(P_0) = f_{10}^{'}$, and $f_2^{'}(0) = f_{20}^{'}$ and substitute (16) and (17) into (15) to obtain

$$(u_1)_1 + (u_a)_1 \approx -(\mu_1(a_1) + (f_{10})_1)u - (f_{10}^1)_1(\rho_0)_1(N_1) - (f_{20}^1)_1(\rho_0)_1(Q_1)$$

$$u_1(t, 0) = \rho_1(t, 0)(\rho_0)_1(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t, a) + (\rho_0)_1(a_1)da_1 - (\rho_0)_1\,(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t, a_1)da_1 + (h_1(Q_1) - 1)(\rho_0)_1\,(0) \qquad (18)$$

Expanding

$$h_1(Q_1) = 1 + (h^1)_1(0)Q_1 + 0(Q_1) \equiv 1 + h_0' Q_1 + 0Q_1$$

And substituting it into (18), we arrive at

$$u_1(t,0) \approx \int_0^\infty \beta_1(a_1) u_1(t,a_1) da_1 + (h_{a_1}')(\rho_0)_1 Q_1$$

The linearization is the following:

$$[u_1] + [u_1] = -(\mu_1(a_1) + [t_{10}]_1 u_1 - [\alpha_1]_1(a_1) N_1 - [\alpha_2]_1(a_1) Q_1$$

$$u_1(t,0) = \int_0^\infty \beta_1(a_1) u_1(t,a_1) da_1 + [\xi_1] Q_1 \qquad (19)$$

$$\frac{dQ_1}{dt} = -b_1 Q_1$$

where

$$[\alpha_1]_1(a_1) = [t_1^1 a]_1 [p,o]_1(a_1) d2(a_1) = [f_{20}^1]_1(\rho\, o\,)_1(a_1)\xi = [\mathrm{h\ o}^1]_1[\rho \mathrm{o}]_1(\mathrm{o}) \text{ and } \mathrm{b}_1 = \mathrm{b}_1 - g_1(P_1)$$

If $g_1(P_0) > b$, clearly, $(\rho_0)_1(a_1), 0)$ is unstable and the predator population is persistent. Assume $g_1(P_0)_1 < b$.

Then, $Q_1(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_1 + u_a = -(\mu_1(a_1) + (f_{10})_1) u_1 - (\alpha_1)_1(a_1) N_1$$

$$= -(\mu_1(a_1) + f_{10}) u_1 - \alpha_1(a_1) \int_0^\infty u_1(t,a_1) da_1$$

$$u_1(t,0) = \int_0^\infty \beta_1(a_1) u_1(t,a_1) da_1$$

Using the method developed by Cushing [2, 3], we suppose

$$u_1(t, a_1) = \tilde{u}_1(a_1)^{e^{\gamma(t-a_1)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_1}\, da_1$. Then,

A solution of the first differential equation for $\tilde{u}(a_1)$ gives,

$$\tilde{u}_1(a_1) = e^{-\int_0^a \mu_1(\tau_1)d\tau - (f_{10})_1 a}\left(\tilde{u}_1(0) - w\int_0^a e^{\int_0^\tau \mu_1(s_1)ds + (f_{10})\tau_1}(\alpha_1)_1(\tau_1)e^{\gamma_1\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_1(a_1)e^{-\gamma_1 a_1 - (f_{10})_1 a_1 - \int_0^a \mu_1(\tau_1)d\tau}\right)\tilde{u}_1(0)$$

$$+w\int_0^\infty \beta_1(a_1)e^{-\gamma_1 a_1 - (f_{10})_1 a_1 - \int_0^a \mu_1(\tau_1)d\tau}\int_0^a (\alpha_1)_1(\tau_1)e^{-\gamma_1\tau_1}e^{(f_{10})_1\tau_1 + \int_0^\tau \mu_1(s_1)ds}d\tau_1\, da_1$$

$$= 0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-ya - f_{10}a - \int_0^a \mu(\tau)d\tau}da\right)u(0)$$

$$+\left(-1 - \int_0^\infty e^{-ya - f_{10}a - \int_0^a \mu(\tau)d\tau}\int_0^a e^{f_{10}\tau + \int_0^\tau \mu(s)ds + y\tau}\alpha_1(\tau)d\tau\, da\right)w = 0. \quad (23)$$

Define $Y(\gamma, a) = e^{-ya - f_{10}a - \int_0^a \mu(\tau)d\tau}$. Then

$$\left(1 - \int_0^\infty \beta\,(a)Y(\gamma, a)da\right)u(0)$$

$+w\int_0^\infty \beta(a)Y(\gamma,a)\int_0^a \alpha_1(\tau)Y^1(\gamma,\tau)d\tau\, da = 0, (24)$

$$\left(\int_0^\infty Y(\gamma,a)da\right)u(0)$$

$$+\left(-1-\int_0^\infty Y(\gamma,a)\int_0^a \alpha_1(\tau)Y^{-1}(\gamma,\tau)d\tau\, da\right)w = 0.$$

Equation (20) has a solution of the form of $u(t,a) = u(a)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $u(0) \not\equiv 0$ and $\gamma, w$ in the complwx plane.

Let $I(\gamma,a) = \int_0^a \rho_0(\tau)Y^{-1}(\gamma,\tau)d\tau$. We obtain the characteristic equation

$$f_{10}^l\left(\int_0^\infty Y(\gamma,a)da\right)\left(\int_0^\infty \beta(a)Y(\gamma,a)I(\gamma,a)da\right)$$

$$+\left(1-\int_0^\infty \beta(a)Y(\gamma,a)da\right)\left(1+f_{10}^l\int_0^\infty Y(\gamma,a)I(\gamma,a)da\right) = 0. \quad (25)$$

Since

$$\rho_0(a) = \rho_0(0)e^{-\int_0^a \mu(\tau)d\tau - f_{10}a}$$

$$I(\gamma,a) = \int_0^a \rho_0(\tau)e^{\gamma\tau + f_{10}\tau + \int_0^\tau \mu(s)ds}d\tau$$

$$= \frac{\rho_0(0)}{\gamma}(e^{\gamma a} - 1)$$

Let

$$E(a) = e^{-f_{10}a - \int_0^a \mu(\tau)d\tau}$$

Then $[Y_1]_1[\gamma_1]_1, a_1) = E_1(a_1)e^{-\gamma_1 a_1}$. by a simple calculation,

$$\int_0^\infty \beta_1(a_1)Y_1(\gamma_1, a_1)I(\gamma_1, a_1)da = \frac{[\rho_0]_1(0)}{\gamma}\left(1 - \int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}\ da\right);$$

$$\int_0^\infty Y_1(\gamma_1, a_1)I(\gamma_1, a_1)da = \frac{[\rho_0]_1(0)}{\gamma_1}\left(\int_0^\infty E_1(a_1)e^{-\gamma_1 a_1}\ da_1\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_1(0)}{\gamma_1}[f_{10}^1]\int_0^\infty E_1(a_1)e^{-\gamma_1 a_1}da_1\ (\int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1)$$

$$+\left(1 - \int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right).$$

$$\left(1 + \frac{(\rho_0)_1(0)}{\gamma_1}(f_{10}^{'})_1\left(\int_0^\infty E_1(\ a_1)da_1 - \int_0^\infty E_1(\ a_1)e^{-\gamma_1 a_1}da_1\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1 - \int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right)(1 + \frac{(f_{10}^{'})_1}{\gamma_1}(\rho_0)_1(0)\int_0^\infty E_1(\ a_1)da_1) = 0.$$

We need ũ(a) $\not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1 - \int_0^\infty \beta_1(a_1)Y_1(\gamma_1 a_1)da_1 = 1 - \int_0^\infty \beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1 \neq 0,$$

Since $\int_0^\infty \beta_1(a_1)Y_1(\gamma_1 a_1)\ \int_0^a \alpha_1\ (\tau_1)Y^{-1}(\gamma_1,\tau_1)d\tau_1 da_1 > 0.$ Therefore, the characteristic solution is

$$\gamma_1 = -(f_{10}^{'})_1(\rho_0)_1 \int_0^\infty E_1(a_1)da_1.$$

Since $f_1'(P_0)_1 > 0,$ we obtain the following result

THEOREM 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_1), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_1), 0)$ which is related to the extinction of debit summation as has been referred to in the foregoing and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(\ p^*(a), Q^*\ )$. It must satisfy the equations:

$$[p_a^*]_1 + (\mu_1(a_1) + [f_1]_1(\ P_1^*) + [f_2]_1(Q_1^*)\ )P_1^* = 0, \quad (26)$$

$$p_1^*(0) = h_1(Q_1^*)\int_0^\infty\ \beta_1(a_1)p_1^*(a_J)p_1^*(a_1)da_1 \quad (27)$$

$$g_1(P_1^*) - b_1 = 0, \quad (28)$$

Where $P_1^*$=$\int_0^\infty p_1^*(a_1)\, da_1$

Equation (28) yields $P^* = g^{-1}(b),$

From (26),

$$p_1^*(a) = p_1^*(0)e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}$$

Substituting into (27), we need to solve

$$\left(1 - h_1(Q_1^*)\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}da_1\right)p_1^*(0) = 0$$

For $Q_1^*$

Fix $p_1^*(0)$ and $P_1^*$ Let

$$H_1(z_1) = 1 - h_1(z_1)\int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1}da_{,1}$$

And assume that

$$[k_3]_1\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(g_1^{-1}(b_1))a_1 - [k_2]_1 a_1}\, da_1 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_1(z_1)] > 0.$$

Clearly,

$$\frac{dH_1}{dz_1} = h_1(z)([f_2']_1(z_1))\int_0^\infty a_1\beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - f_1(P_1^*)a_1 - f_2(z_1)a_1}da_{\square}$$

$-h'(z_1)\int_0^\infty \beta_1(a_1)\, e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1}\quad da_1$

By the assumptions on $\left[f_2^1\right]_1 and\ [h^1]^1 dH_1/dz > 0$. Hence, if

$$H_1(0) = 1 - \int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1} da_1 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_1(0) < 0$ follows if $g_1^-(b_1) < [f_1^{-1}]_1(r_1)$, where $r_1$ is the solution of (12).

Indeed, since

$$R_1(s_1) = \int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 e^{-y}1^a} da_1$$

Is a decreasing function of $s_1$ when $P_1^* = [g^{-1}]_1(b_1) < [f_1^{-1}]_1(r_1)$, that is $[f_1]_1(P_1^*) < r$, $R\left(f_1(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_1(g^{-1}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1\left(g_1^{-1}(b)\right)a_1} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_1^*(a_1), Q_1^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_1^*(a_1), Q_1^*)$ is more complicated. We again consider

perturbations $m_1(t_1 a_1) = p_1(t_1, a_1) - p_1^*(a_1)$ and $W_1(t) = Q_1(t) - Q^*$.Then

$$M_1(t) \equiv \int_0^\infty m_1(t, a_1) da_1 = P_1(t) - P_1^*,$$

$$[m_1]_1 + [m_a]_1 = [p_1]_1 + [p_a]_1 - \left[p^*\right]_1(a_1) \approx -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1$$

$$-[f_{1*}^{1}]_1 p_1^*(a_1)M - [f_{2*}^{1}]_1 p_1^*(\mathrm{a})W_1,$$

$$m_1(t,0)p_1(t,0) - p_1^*(0) \approx h_* \int_0^\infty \beta_1(a_1)m_1(t,a_1)\,da + \frac{[h_*^1]_1}{[h_*]_1}\, p_1^*(0)W_1$$

$$\frac{dW_1}{dt} = \frac{dQ_1}{dt} \approx [g_*^1]_1 Q_1^* M_{,1}$$

Where $[f_{1*}]_1$ denoting $[f_1(P^*)]_1$, $[f_{1*}]_1$ denoting $[f_2(P^*)]_1$, $\left[f_{1*}^{'}\right]_1$ denoting $\left[f^{'}{}_1(P^*)\right]_1 \left[f_{2*}^{'}\right]_1$ denoting

$[f_2]_1$ $(Q_1^*,$ $h_{1}^{*}$ denoting $h_*^{'}$ denoting $h_1^{'}$ $(Q_1^*)$, and $\left[g_*^{'}\right]_1$

Denoting $\left[g^{'}\right]_1 (P_1^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_1]_1 + [m_a]_1 = -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1 - \left[f^{'}{}_{1*}\right]_1 [p_1^*](a_1)M_1 - \left[f_{2*}^{'}\right]_1 p_1^*(a_1)W_1$$

$$m_1(t,0) = [h_*]_1 \int_0^\infty \beta_1(a_1)m_1(t,a_1)da_1 + \frac{[h_*^1]_1}{[h^*]_1}[p^*]_1(0)\, W_{,1} \qquad (30)$$

$$\frac{dW_1}{dt} = [g_*^1]_1 [Q_1^*] M_1$$

Integrating $[m_1]_1(t,a_1)$ with respect to a from 0 to ∞, we have

$$\frac{dM_1(t)}{dt} = \int_0^\infty [mt]_1(t,a_1)da_1$$

$$= -\int_0^\infty [m_a]_1\,(t,a_1)da_1 - \int_0^\infty \mu(a_1)m(t,a_1)da_1$$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f'_{1*}]_1 \int_0^\infty p_1^*(a_1)da_1\right) M_1 - [f_{2*}^1]_1 \int_0^\infty p_1^*(a_1)da_1 W_1$$

$= m_1(t,0) - \int_0^\infty \mu_1(a_1)\, m_1(t,a_1)da_1$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f_{1*}^{i}]_1 P_1^*\right) M_1 - [f_{2*}^{i}]_1 P_1^* W_{\cdot 1}$$

From (30), we arrive at the equations

$$\frac{dW_1(t)}{dt} = \int_0^\infty ([h_*]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1$$

$$-\left([f_{1*}] + [f_{2*}]_1 + [f_{1*}']_1 P_1^*\right) M_1 + \left(\frac{[h_*']_1}{[h_*]^1}\, P_1^*(0) - f_{2*}' P_1^*\right) W_1,$$

(31)

$$\frac{dW_1(t)}{dt}[g'_*]_1 Q_1^*\, M_{\cdot 1}$$

By the mean value theorem,

$$\int_0^\infty ([h_{"*}]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1 = ([h_*]_1\beta_1(\theta_1) - \mu_1(\theta_1))M_{,1}$$

Where $0< \theta_1 = \theta_1(t) < \infty$.

And if $a_1^2 [a_1^2]_1 < 4 [a_2]_1 < 4 a_2 e$,

Re $\lambda_1[\lambda_1]_1 = \mathrm{Re}\lambda_2[\lambda_2]_1 = -\frac{1}{2}a_1[a2\ ]_1 < 0.$

Let $\Delta$=max$_{2<u<\infty}\left|h_*[h_*]_1\beta_1(a_1) - \mu_1(a_1)\right|$. Then$|n_1(t)|$ ⊿∀t ⓪. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium($\rho^*(a), \boldsymbol{Q}^*$) is asymptotically stable if either $a_1^2 < 4\,a_2 e$ and $[a_1] > 2\Delta$,or $[a_1^2]$ ④$[a_2]$e and $(a_1 - \sqrt{(a_1^2)_1 - 4(a_2)_1 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+1}\right)$ and $\beta_1(a_1) = ae^{-ma}$ Then $\Delta$=$[\mu_0] + 1$ Simple calculation leads to

$$\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(\tau_1) + d\tau_1 da} = \frac{1}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_1 + [\mu_0]$ ⓪.5. Assume$[f_1](x_1)$= $K_1x_1$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_1(x_1)$= $[0x_1]/(1 + x_1)$; the functional responses $[[f_2](x) = [0]\, x_1/(1 + x_1)$and $h_1(x_1)$=$(1+\omega_1x_1)/(1 + x_1)$(See Holling [7] and Rashevsky [13])

Equation (14),now has the form of

$$\frac{m_1 + [\mu_0] + [\mathrm{K}][P_0])}{[m_1]_1 + (\mu_0)_1 + [\mathrm{K}][P_0])_1^3\ [\mathrm{K}]_1[P_0]_1^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1+P_0) < b_1$ i e $[P_0] < b/(\delta_1 - b_1)$then the equilibrium$(P_0)$ $[P_0]_1(a), 0)$ is stable This implies that the

cedit confederation goes to extinction but that the cedit confederation persists $b_1$

$[f_{1*}] = K_1$ P $^*$ $[f_{2*}] = \frac{\sigma Q_1^*}{1+Q_1^*}$ $[h_*] = \frac{1+\omega_1 Q_1^*}{1+Q_1^*}$

$[f_{1*}] = K_1$ $\left[f_{2*}^{1}\right] = \frac{\sigma_1}{(1+Q_1^*)^2}$ $[g_*^1] = \frac{\sigma_1}{(1+P_1^*)^2}$ $\left[h_*^1\right] = \frac{\omega_1 - 1}{(1+Q_1^*)^2}$

## AN ANALOGY BETWEEN PREDATOR PREY AND ACCENTUATION ATTRION MODELS

**Best example how the model could be applied is the Schrödinger's equation itself wherein RHS gobbles up LHS, and the terms on RHS do the same with each other. It is like adding milk to water. Whether it happens instantaneously or time lag is there is a different question altogether dealt at the opportune points**

**Best example that can be given is the formulation of the quantum wave equation albeit in a nascent form for the recently found gravitational waves. This has been studied in length in various papers of gravitational waves, wherein the attribution of how the model is used is described. Note the calculation holds for only one module. This is in consideration to the fact that the relation is beteen the both debit systems and credit systems which as we have enucleated in the foregoing bears ample testimony for the application of the model**

**Gravitational quantum wave equation:**

**Equation Time-dependent equation (Wikipedia)**

**(1) In quantum mechanics, the Schrödinger equation is (=) a partial differential equation that describes how the quantum state of a quantum system changes with (e&eb) time. It was formulated in late 1925, and published in 1926, by the Austrian physicist Erwin Schrödinger.[1]**

**(2) In classical mechanics Newton's second law, (F = ma), is used to make (eb) a mathematical prediction as to what path a given system will take following a set of known initial conditions.**

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| **(3) In quantum mechanics, the analogue of Newton's law is (=) Schrödinger's equation for (e) a quantum system (usually atoms, molecules, and subatomic particles whether free, bound, or localised).** |
| **(4) It is not a simple algebraic equation, but in general a linear partial differential equation, describing (eb) the time-evolution of the system's wave function (also called a "state function").[2]:1–2** |
| **(5) The concept of a wavefunction is (=) a fundamental postulate of quantum mechanics.** |
| **(6) Although Schrödinger's equation is often presented as a separate postulate, some authors [3]: Chapter 3 show that some properties resulting from Schrödinger's equation may be deduced just from (e) symmetry principles alone, for example the commutation relations.** |
| **(7) Generally, "derivations" of the Schrödinger equation demonstrate (eb) its mathematical plausibility for describing (eb) wave–particle duality, but to date there are no universally accepted derivations of (e) Schrödinger's equation from (e) appropriate axioms.** |
| **(8) In the Copenhagen interpretation of quantum mechanics, the wave function is (=) the most complete description that can be given of (e) a physical system.** |
| **(9) Solutions to Schrödinger's equation describe (eb) not only molecular, atomic, and subatomic systems, but also macroscopic systems, possibly even the whole universe.[4]:292ff** |
| **(10) The Schrödinger equation is consistent with (e&eb) classical mechanics but not with (e&eb) special relativity.** |
| **(11) Making quantum mechanics consistent with (e&eb) special relativity is one of the great achievements of (e) Quantum Field Theory and requires (e) the introduction of the creation and annihilation operators so that (e) the number of particles can no longer be considered constant.** |
| **(12) The Schrödinger equation is not the only way to make predictions in (eb) quantum mechanics—other formulations can be used, such as Werner Heisenberg's matrix mechanics, and Richard Feynman's path integral formulation** |
| **(13) The form of the Schrödinger equation depends on (=) the physical situation (see below for special cases).** |
| **(14) The most general form is the time-dependent Schrödinger equation, which gives a description of a system evolving with time: [5]:143** |

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| **(15) A wave function that satisfies the non-relativistic Schrödinger equation with V = 0. In other words, this corresponds to (e &eb) a particle traveling freely through (e&eb) empty space. Time-dependent** |
| **Schrödinger equation (general)** |
| **(16) i ℏ ∂ ∂ t Ψ ( r , t ) = H ^ Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t) where i is the imaginary unit, ħ is the Planck constant divided by 2π, the symbol ∂/∂t indicates a partial derivative with respect to time t, Ψ (the Greek letter psi) is the wave function of (e) the quantum system, and Ĥ is the Hamiltonian operator (which characterizes the total energy of any given wave function and takes different forms depending on the situation).** |
| **(17) Each of these three rows is a wave function which satisfies (eb) the time-dependent Schrödinger equation for a harmonic oscillator. Left: The real part (blue) and imaginary part (red) of the wave function. Right:** |
| **(18) The probability distribution of finding the particle with (e&eb) this wave function at a given position.** |
| **(19) The top two rows are examples of stationary states, which correspond to (e&eb) standing waves.** |
| **(20) The bottom row is an example of a state which is (=) not a stationary state.** |
| **(21) The right column illustrates (eb) why stationary states are (=) called "stationary".** |
| **(22) The most famous example is the non-relativistic Schrödinger equation for (e) a single particle moving in an electric field (but not a magnetic field; see the Pauli equation): [6]** |
| **Time-dependent Schrödinger equation** |
| **(Single non-relativistic particle)** |
| **(23) i ℏ ∂ ∂ t Ψ ( r , t ) = [ − ℏ 2 2 μ ∇ 2 + V ( r , t ) ] Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)** |
| **(24) Here μ is (=) the particle's "reduced mass", V is its potential energy, ∇2 is the Laplacian (a differential operator), and Ψ is (=) the wave function (more precisely, in this context, it is called the (=) "position-space wave function").** |

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| **(25) In plain language, it means "total energy equals (=) kinetic energy plus potential energy", but the terms take unfamiliar forms for (e) reasons explained below.** |
| **(26) Given the particular differential operators involved, this is (=) a linear partial differential equation.** |
| **(27) It is also a diffusion equation, but unlike the heat equation, this one is (=) also a wave equation given (e&eb) the imaginary unit present in the transient term.** |
| **(28) The term "Schrödinger equation" can refer to (e) both the general equation (first box above), or the specific nonrelativistic version (second box above and variations thereof).** |
| **(29) The general equation is indeed quite general, used throughout quantum mechanics, for everything from the Dirac equation to quantum field theory, by plugging in various complicated expressions for the Hamiltonian. The specific nonrelativistic version is a simplified approximation to reality, which is quite accurate in many situations, but very inaccurate in others (see relativistic quantum mechanics and relativistic quantum field theory)** |
| **(30) To apply the Schrödinger equation, the Hamiltonian operator is set up for the system, accounting for the kinetic and potential energy of the particles constituting the system, then inserted into the Schrödinger equation. The resulting partial differential equation is solved for the wave function, which contains information about the system.** |
| **Time-independent equation** |
| **(31) The time-dependent Schrödinger equation described above predicts that wave functions can form standing waves, called stationary states (also called "orbitals", as in atomic orbitals or molecular orbitals). These states are important in their own right, and if the stationary states are classified and understood, then it becomes easier to solve the time-dependent Schrödinger equation for any state. Stationary states can also be described by a simpler form of the Schrödinger equation, the time-independent Schrödinger equation. (This is only used when the Hamiltonian itself is not dependent on time explicitly. However, even in this case the total wave function still has a time dependency.)** |
| **Time-independent Schrödinger equation (general)** |
| **(32) E Ψ = H ^ ⊠ Ψ {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi } {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi }** |
| **In words, the equation states:** |
| **(33) When the Hamiltonian operator acts on a certain wave function Ψ, and the result is proportional to the same wave function Ψ, then Ψ is a stationary state, and the proportionality constant, E, is the energy of the state Ψ.** |

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| **(34) The time-independent Schrödinger equation is discussed further below. In linear algebra terminology, this equation is an eigenvalue equation.** |
| **(35) As before, the most famous manifestation is the non-relativistic Schrödinger equation for a single particle moving in an electric field (but not a magnetic field):** |
| **Time-independent Schrödinger equation (single non-relativistic particle)** |
| **(36) E Ψ ( r ) = [ − ℏ 2 2 μ ∇ 2 + V ( r ) ] Ψ ( r ) {\displaystyle E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )} E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )** |
| **Gravitational waves: (Wikipedia)** |
| **(37) In theory, the loss of energy through gravitational radiation could eventually drop the Earth into the Sun. However, the total energy of the Earth orbiting the Sun (kinetic energy + gravitational potential energy) is about 1.14×1036 joules of which only 200 Watts (joules per second) is lost through gravitational radiation, leading to a decay in the orbit by about 1×10−15 meters per day or roughly the diameter of a proton. At this rate, it would take the Earth approximately 1×1013 times more than the current age of the Universe to spiral onto the Sun. This estimate overlooks the decrease in r over time, but the majority of the time the bodies are far apart and only radiating slowly, so the difference is unimportant in this example.** |
| **More generally, the rate of orbital decay can be approximated by[39]** |
| **(38) d r d t = − 64 5 G 3 c 5 ( m 1 m 2 ) ( m 1 + m 2 ) r 3 , {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,} {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,}** |
| **where r is the separation between the bodies, t time, G Newton's constant, c the speed of light, and m1 and m2 the masses of the bodies. This leads to an expected time to merger of [39]** |
| **t = 5 256 c 5 G 3 r 4 ( m 1 m 2 ) ( m 1 + m 2 ) . {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.} {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.}** |
| **(39) For example, a pair of solar mass neutron stars in a circular orbit at a separation of 1.89×108 m (189,000 km) has an orbital period of 1,000 seconds, and an expected lifetime of 1.30×1013 seconds or about 414,000 years. Such a system could be observed by LISA if it were not too far away. A far greater number of white dwarf binaries exist with orbital periods in this range. White dwarf binaries have masses in the order of the Sun, and** |

**diameters in the order of the Earth. They cannot get much closer together than 10,000 km before they will merge and explode in a supernova which would also end the emission of gravitational waves. Until then, their gravitational radiation would be comparable to that of a neutron star binary.**

**(40) By substituting value of dr/dt in Schrödinger's wave equation we get the quantum wave equation for gravitational waves.**

In order to have (H3) and (H4) hold, $m, \mu_0, \mathrm{K}, \mathrm{b}, \delta$, and $\omega$ need to satisfy

$$\frac{\omega_1}{(m+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1+b_1}+\sigma)^2}+\frac{2\omega_1}{([m_1]_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1+b_1}+\sigma_1)^3}<1$$

And

$$\frac{1}{(m_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1-b_1})^2}+\frac{2\omega_1}{(m_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1-b_1})^3}>1$$

When the inequalities (33) and (34) hold there is a unique equilibrium $\mathrm{Q}^*$ which can be solved from the following equation:

$$(m_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1-b_1}+(m_1+[\mu_0]_1+\frac{K_1 b_1}{\delta_1-b_1}+\sigma_1)\,\mathrm{Q}_1^*)^3$$

$$=m_1+[\mu_0]_1+3+\frac{K_1 b_1}{\delta_1-b_1}(1+\mathrm{Q}^*)(1+\omega_1\mathrm{Q}_1^*$$

Now

$$\left[\mathrm{a}_1^2\right]-4\left[\mathrm{a}_2\right]\mathrm{e}=\left(\frac{2K_1 b_1}{\delta_1-b_1}\right)^2+\frac{4K_1 b_1\sigma_1\,\mathrm{Q}_1^*}{(\delta_1-b_1)(1+\mathrm{Q}_1^*)}+\frac{\sigma_1^2\,\mathrm{Q}_1^{*2}}{(1+\mathrm{Q}_1^*)^2}$$

$$\frac{4\,b_1\sigma_1\,(\delta_1-\mathrm{b})\mathrm{Q}^*}{\delta_1\,(1+\mathrm{Q}_1^*)^2}-\frac{4\,(\delta_1-b_1)^2(1-\omega_1)\mathrm{Q}_1^*}{\delta_1\,(1+\omega_1\mathrm{Q}_1^*)(1+\mathrm{Q}_1^*)}\mathrm{P}^*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_1\,ae^{mu}_1\;\;[\mu_0]_1-\frac{a_1}{a_1+1}$$

Hence if $[\mathrm{a}_1^2>4(a_2)_1\mathrm{e}$ and

$$[a_1]=\frac{2K_1b_1}{\delta_1-b_1}+\frac{\sigma_1\,\mathrm{Q}_1^*}{1+\mathrm{Q}_1^*}>2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_1]_1=[f_{1*}]_1+[f_{2*}]_1+[f_{1*}]\,\mathrm{P}_1^*$$

$$[a_2]_1=\left[\mathrm{f}_{2*}^{1}\right]\mathrm{P}_1^*-\frac{\mathrm{h}_*^1}{\mathrm{h}_*}\mathrm{P}^*(0)$$

$$\mathrm{e}=[\,\mathrm{g}_*^1]_1\,\mathrm{Q}_1^*$$

$$\eta_1(\mathrm{t})=\left[\mathrm{h}_*\right]_1\beta_1\left(\theta_1(\mathrm{t})\right)-\mu_1(\theta(\mathrm{t})),$$

The system(31)can be written as the following system:

$$[\mathrm{M(t)}]\left[\mathrm{a}_1\right]M_1-[a_2]_1W_1-[\mathrm{n(t)}]M_1$$

$$w_1(t) = eM_1$$

Write

$$A=\begin{pmatrix} -[a_1]_1 & -[a_2]_1 \\ e & 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_1 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_1 \\ W_1 \end{pmatrix}$$

The solution of (32)is

$$X_1(t) = e^{A,t}[X]_1(0) + \int_0 e^{At}$$

## Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

## (Module Three)

**Abstract: Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-435) 1990. We demonstrate its applicability cosmology/ quantum computation economics et al. As stated in the foregoing we shall apply model to various cosmological problem and show the relevance to the same. . To exemplify model can be applied to the following situational significations: Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new**

**source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 Masaru Shibata and Takashi Nakamura** perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. **Authors** use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society **Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995**

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(g_3)Q_3$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_3(0) = 0, \quad g_3(P_3) > 0, \quad P_3 \; 0, \qquad \lim_{P_3 \to \infty} g_3(P_3) = (K_3)_3 \; \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_3 + (\rho_a)_3 + D_3(a_3, P_3, Q_3)\rho_3 + 0, \qquad 0 < a_3 \; \infty,$$

$$\rho_3(t,0) = \int_0^\infty n_3(a_3, P_3, Q_3)\rho_3(t_3, a_3)da_3$$

Where $D_3$ is the mortality function and $n_3$ is the fertility function.

Put $_3$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_3}{dt}\right]_3 = x_3\left[\frac{dQ_2}{dt}\right]_3 - x_2\left[\frac{dQ_3}{dt}\right]_3$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_3 - x_4\left[\frac{dQ_2}{dt}\right]_3$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_3 - x_6\left[\frac{dQ_3}{dt}\right]_3$$

Where $x_3, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_3$ and $g_3(P_3)Q_3$

Now when you add the term $(b_3Q_3)$ with $g_3(P_3)(Q_3)$ we are infect see is that $(b_3Q_3)$ is credited incrementally part of $g_3(P_3)(Q_3)$ or instant nearly, when subtracted from $\frac{dQ_3}{dt}$ , it means that $g_3(P_3)(Q_3)$ is slowly eaten by $\frac{dQ_3}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_3 = \mu_3(a_3) + (f_3)_3(P_3)_3 + (f_2)_3(Q)_3 \quad (2)$$

$$n_3 = \beta_3(a_3)h_3(Q_3), \quad (3)$$

Where $\mu_3$the age- dependent debit rate whit is $\mu_3(a_3)$ $\forall a > 0$; $(f_3)_3$ is the density- dependent death rate with $f_3(0) = 0$, $(f_3^3)_3(P_3)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_3)_3(P_3) = (k_3)_3$ ; $f_2$ is the functional response of creditors to prey with $[f_2]_3(0) = 0, (f_2^3)_3(Q) > 0, \forall Q$ , and $\lim_{Q\to\infty}(f_2)_3(Q_3) = [k_2]$ ; $\beta$ is the age dependent birth rate with $\beta(a)$ $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 3, h^3(Q)$ $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ .

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_3 + (\rho_a)_3 = -\mu_3(a_3) + (f_3)_3P_3) + \big((f_2)_3(Q_3)\big)(\rho_3)_3 \quad (4)$$

$$\rho_3(t_3, 0) = h_3(Q_3)\int_0^\infty \beta_3(a_3)\rho_3(t, a_3)da, \quad (5)$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(P_3)Q_3 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_3$ and $Q_3$ of (4-6) satisfy

$$0 \le \rho_3(t, a_3) \le \rho_3^*(t, a_3),$$

$$0 \le Q_3(t) \le Q_3^*(t),$$

Where $\rho_3^*(t, a_3)$ and $Q_3^*(t)$, are solutions of

$$(\rho_t^*)_3 + (\rho_a^*)_3 = -(\mu_3(a_3) + (f_3)_3(P_3^*) + (f_2)_3 Q_3^*))\,\rho_3^*, \quad (7)$$

$$\rho_3^*(t, 0) = \int_0^\infty \beta_3(a_3)\,\rho_3^*(t, a_3)da_3 \quad (8)$$

$$\frac{dQ_3^*}{dt} = -b_3 Q_3^* + g_3(P_3^*)Q_3^*, \quad (9)$$

With $P_3^*(t) = \int_0^\infty \rho_3^*(t, a_3)da_3$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_3^*(t)}{dt} = P_3^*(t)\big(r_3 - (f_3)(P_3^*) - (f_2)_3(Q^*)\big), \quad (10)$$

$$\frac{dQ_3^*(t)}{dt} = Q_3^*(t)(-b_3 + g_3(P_3^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_3(a)e^{-\int_0^a \mu_3(\tau)dt}e^{-ra}\,da_3 = 1. \qquad (12)$$

If $r < 0$, the trivial solution $P_3^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(t)dt}\,da_3 > 1, \qquad (H1)$$

Which guarantees $r_3 > 0$.

*Remark.* If $\beta_3(a) = [\beta_0]_3 a_3 e^{-\alpha_3 a_3}, \mu_3(a_3) \equiv \mu_3$ then the solution of (12) is

$$r_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_3]_3(P_3) = 0, [f_2](Q_3) = \lambda_3 Q_3\ h_3(Q_3) = 1, \qquad and\ \ g_3(P_3) = c_3 P_3$$

$$[\rho_3]_3 + [\rho_a]_3 = -(\mu_3 + \lambda_3 Q_3)\rho_3$$

$$\rho_3(t,0) = \int_0^\infty [\beta_0]_3 a_3 e^{-\alpha_3 a_3}\rho_3(t,a_3)\,da_3$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + c_3 P_3 Q_3$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_3}{dt} = \omega_3 P_3 - \lambda_3 P_3 Q_3$$

$$\frac{dQ_3}{dt} = b_3 Q_3 - c_3 P_3 Q_3$$

Where $\omega_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$

## THE EQUILIBRIUM $(\rho_0)_3(a_3), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_3(a_3), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_3 + (\rho_a)_3 = -\big(\mu_3(a_3) + (f_3)_3(P_3)\big)\rho,$$

$$\rho_3(t,0) = \int_0^\infty \beta_3(a_3)\rho_3(t,a_3)da_3 \qquad (13)$$

The solution of (13) has the form of

$$\rho_3(t,a_3) = Ce_3^{r(t-a_3)} - \int_0^a (\tau)d\tau - \int_0^3 (f_3)_3\big(P_3(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_3(t) = C_3 e^{r_3 t - \int_0^t (f_3)_3(P_3(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_3^{-3}(r)$ if $k_3 > r$. This is equivalent to

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{(k_3)_3 a_3} da_3 < 1.$$

In fact, if we define

$$R_3(s_3) = \int_0^{\infty} \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau)d\tau} e^{-s_3 a_3}\, da_3$$

$R(s)$ is a decreasing function. $R_3(r_3) = 1$ And $R_3(k_3)_3 < 1$ imply $r_3 < (k_3)_3$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_3(a_3), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{k_3 a_3} da_3$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_3 \equiv \int_0^{\infty}(\rho_0)_{33}(a_3)da_3$ *satisfies*

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{-(f_3)_3(P_0)a_3} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_3(a_3), 0)$

Let $u_3(t, a_3) = \rho(t, a_3) - \rho_0(a_3)$ Then

$$N_3(t) \equiv \int_0^{\infty} u_3\,(t, a_3)da_3 = \int_0^{\infty} \rho_3(t, a_3)da_3 - \int_0^{\infty}(\rho_0)_3(a_3)da = P_3(t) - P_0)_3$$

And

$$(u_3)_3 + (u_a)_3 = (\rho_t)_3 + (\rho_a)_3 - (\rho_0^3)_3(a_3)$$

$$= -\mu_3(a_3) + (f_3)_3(P_3) + (f_2)_3(Q_3))\,(u_3 + (\rho_0)_3) + \mu_3(a_3)(\rho_0)_3 + (f_3)_3(P_0)_3(\rho_0)_3$$

$$= -\mu_3(a_3)u_3 - (f_3)_3(P_3) + (f_2)_3(Q_3))\,u_3 - (\rho_0)_3\,(f_3)_3(P_3) - (f_3)_3(P_0)_3$$

$$+ (f_2)_3(Q_3)), \qquad (15)$$

Where $(\rho_0^3)_3(a_3)$ denotes $d(\rho_0)_3(a_3)/da_3$

Expansion $(f_3)_3(P_3)$ and $(f_2)_3(Q_3)$ at $(P_0)_3$ and 0 yield respectively

$$(f_3)_3(P_3) = (f_3)_3((P_0)_3 + N_3)$$

$$= (f_3)_3(P_0)_3 + (f_3^3)_3(P_0)_3N_3 + o_3(N_3), \qquad (16)$$

$$(f_2)_3(Q_3)$$

$$= (f_2^3)_3(0)Q + o_3(Q_3) \qquad (17)$$

Write $f_3(P_0) = f_{30}, f_3'(P_0) = f_{30}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_3)_3 + (u_a)_3 \approx -(\mu_3(a_3) + (f_{30})_3)u - (f_{30}^3)_3(\rho_0)_3(N_3) - (f_{20}^3)_3(\rho_0)_3(Q_3)$$

$$u_3(t,0) = \rho_3(t,0)(\rho_0)_3(0)$$

$$= h_3(Q_3)\int_0^\infty \beta_3(a_3)(u_3(t,a) + (\rho_0)_3(a_3)da_3 - (\rho_0)_3\,(0)$$

$$= h_3(Q_3)\int_0^\infty \beta_3(a_3)(u_3(t,a_3)da_3 + (h_3(Q_3) - 3)(\rho_0)_3\,(0) \qquad (18)$$

Expanding

$$h_3(Q_3) = 1 + (h^3)_3(0)Q_3 + 0(Q_3) \equiv 1 + h_0'Q_3 + 0Q_3$$

And substituting it into (38), we arrive at

$$u_3(t,0) \approx \int_0^\infty \beta_3\,(a_3)u_3(t,a_3)da_3 + (h_{a_3}')(\rho_0)_3Q_3$$

The linearization is the following :

$$[u_3] + [u_3] = -(\mu_3(a_3) + [t_{30}]_3 u_3 - [\alpha_3]_3(a_3)N_3 - [\alpha_2]_3(a_3)Q_3$$

$$u_3(t,0) = \int_0^\infty \beta_3(a_3)u_3(t,a_3)da_3 + [\xi_3]Q_3 \qquad (19)$$

$$\frac{dQ_3}{dt} = -b_3 Q_3$$

where

$$[\alpha_3]_3(a_3) = [t_3^3 a]_3[p,o]_3(a_3)d2(a_3) = [f_{20}^3]_3(\rho\, o\,)_3(a_3)\xi = [\mathrm{h\ o^3}]_3[\rho\mathrm{o}]_3(\mathrm{o}) \text{ and } \mathrm{b_3} = \mathrm{b_3} - g_3(P_3)$$

If $g_3(P_0) > b$, clearly, $(\rho_0)_3(a_3), 0)$ is unstable and the predator population is persistent. Assume $g_3(P_0)_3 < b$.

Then, $Q_3(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_3 + u_a = -(\mu_3(a_3) + (f_{30})_3)u_3 - (\alpha_3)_3(a_3)N_3$$

$$= -(\mu_3(a_3) + f_{30})u_3 - \alpha_3(a_3)\int_0^\infty u_3(t,a_3)da_3$$

$$u_3(t,0) = \int_0^\infty \beta_3\,(a_3)u_3(t,a_3)da_3$$

Using the method developed by Cushing [2, 3], we suppose

$$u_3(t,a_3) = \tilde{\mathrm{u}}_3(a_3)^{e^{\gamma(t-a_3)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_3}\, da_3$. Then,

A solution of the first differential equation for $\tilde{u}(a_3)$ gives,

$$\tilde{u}_3(a_3) = e^{-\int_0^a \mu_3(\tau_3)d\tau - (f_{30})_3 a}\left(\tilde{u}_3(0) - w\int_0^a e^{\int_0^\tau \mu_3(s_3)ds + (f_{30})}\tau_3\,(\alpha_3)_3(\tau_3)e^{\gamma_3\tau}d\tau\right).$$

Substituting the expression (13) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\right)\tilde{u}_3(0)$$

$$+w\int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\int_0^a (\alpha_3)_3(\tau_3)e^{-\gamma_3\tau_3}e^{(f_{30})_3\tau_3 + \int_0^\tau \mu_3(s_3)ds}d\tau_3\, da_3$$

$$= 0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}da\right)(\tilde{u}_3(0))_3$$

$$+\left(-3 - \int_0^\infty e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\int_0^a e^{(f_{30})_3\tau_3 + \int_0^\tau \mu_3(s_3)ds + y\tau}\alpha_3(\tau)d\tau\, da\right)w = 0.$$

(23)

Define $Y(\gamma, a) = e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(t_3)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_3\,(a_3)Y_3(\gamma_3,a_3)da_3\right)u_3(0)$$

$$+w\int_0^{\infty}\beta_3(a_3)Y_3(\gamma_3,a_3)\int_0^{a}(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\,da_3=0,$$

(24)

$$\left(\int_0^{\infty}Y_3(\gamma_3,a_3)da_3\right)\tilde{u}_3(0)$$

$$+\left(-1-\int_0^{\infty}Y_3(\gamma_3,a_3)\int_0^{a}(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\,da_3\right)w=0.$$

Equqtion (20) has a solution of the form of $u_3(t,a_3)=u_3(a_3)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_3(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_3(\gamma_3,a_3)=\int_0^{a}(\rho_0)_3(\tau_3)Y^{-3}\,(\gamma_3,a_3)d\tau_3$. We obtain the characteristic equation

$$f_{30_3}^{l}\left(\int_0^{\infty}Y_3(\gamma_3,a_3)da_3\right)\left(\int_0^{\infty}\beta_3(a_3)Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$+\left(1-\int_0^{\infty}\beta_3(a_3)Y_3(\gamma_3,a_3)da\right)\left(1+(f_{30}^{l})_3\int_0^{\infty}Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_3(a_3)=(\rho_0)_3(0)e^{-\int_0^{a}\mu_3(\tau_3)d\tau_3-(f_{30})_3a_3}$$

$$I_3(\gamma_3, a_3) = \int_0^a (\rho_0)_3(\tau_3) e^{\gamma_3\tau_3 + (f_{30})_3\tau + \int_0^\tau \mu_3(s_3)ds_3} d\tau$$

$$= \frac{(\rho_0)_3(0)}{\gamma_3}(e^{\gamma_3 a_3} - 1)$$

Let

$$E_3(a_3) = e^{-(f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau_3}$$

Then $[Y_3]_3[\gamma_3]_3, a_3) = E_3(a_3)e^{-\gamma_3 a_3}$. by a simple calculation,

$$\int_0^\infty \beta_3(a_3)Y_3(\gamma_3, a_3)I(\gamma_3, a_3)da = \frac{[\rho_0]_3(0)}{\gamma}\left(1 - \int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}\ da\right);$$

$$\int_0^\infty Y_3(\gamma_3, a_3)I(\gamma_3, a_3)da = \frac{[\rho_0]_3(0)}{\gamma_3}\left(\int_0^\infty E_3(a_3)e^{-\gamma_3 a_3}\ da_3\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_3(0)}{\gamma_3}[f_{30}^3]\int_0^\infty E_3(a_3)e^{-\gamma_3 a_3}da_3\ (\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3)$$

$$+\left(1 - \int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right).$$

$$\left(1 + \frac{(\rho_0)_3(0)}{\gamma_3}(f_{30}^{'})_3\left(\int_0^\infty E_3(a_3)da_3 - \int_0^\infty E_3(a_3)e^{-\gamma_3 a_3}da_3\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right)(1+\frac{(f_{30}^{'})_3}{\gamma_3}(\rho_0)_3(0)\int_0^{\infty}E_3(a_3)da_3)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_3(a_3)Y_3(\gamma_3 a_3)da_3=1-\int_0^{\infty}\beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\neq 0,$$

Since $\int_0^{\infty}\beta_3(a_3)Y_3(\gamma_3 a_3)\int_0^{a}\alpha_3(\tau_3)Y^{-3}(\gamma_3,\tau_3)d\tau_3 da_3>0$. Therefore, the characteristic solution is

$$\gamma_3=-(f_{30}^{'})_3(\rho_0)_3\int_0^{\infty}E_3(a_3)da_3.$$

Since $f_3'(P_0)_3>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_3),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_3),0)$ which is related to the extinction of debit summation types

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a),Q^*)$. It must satisfy the equations:

$$[p_a^*]_3 + (\mu_3(a_3) + [f_3]_3(P_3^*) + [f_2]_3(Q_3^*))P_3^* = 0, \quad (26)$$

$$p_3^*(0) = h_3(Q_3^*)\int_0^\infty \beta_3(a_3)p_3^*(a_J)p_3^*(a_3)da_3 \quad (27)$$

$$g_3(P_3^*) - b_3 = 0, \quad (28)$$

Where $P_3^* = \int_0^\infty p_3^*(a_3)\, da_3$

Equation (28) yields $P^* = g^{-3}(b)$,

From (26),

$$p_3^*(a) = p_3^*(0)e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}$$

Substituting into (27), we need to solve

$$\left(1 - h_3(Q_3^*)\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}da_3\right)p_3^*(0) = 0$$

For $Q_3^*$

Fix $p_3^*(0)$ and $P_3^*$ Let

$$H_3(z_3) = 1 - h_3(z_3)\int_0^\infty \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3}da_{,3}$$

And assume that

$$[k_3]_3\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(g_3^{-3}(b_3))a_3 - [k_2]_3a_3}\, da_3 < 1. \quad (H3)$$

Then

$$\lim_{z\to\infty}[H_3(z_3)] > 0.$$

Clearly,

$$\frac{dH_3}{dz_3} = h_3(z)([f_2^{'}]_3(z_3)) \int_0^{\infty} a_3\beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - f_3(P_3^*)a_3 - f_2(z_3)a_3}\, da_3$$

$-h^{'}(z_3)\int_0^{\infty} \beta_3(a_3)\, e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3}\ da_3$

By the assumptions on $\left[f_2^3\right]_3 and\ [h^3]^3 dH_3/dz > 0$. Hence, if

$$H_3(0) = 1 - \int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3}\, da_3 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_3(0) < 0$ follows if $g_3^{-}(b_3) < [f_3^{-3}]_3(r_3)$, where $r_3$ is the solution of (12).

Indeed, since

$$R_3(s_3) = \int_0^{\infty} \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 e^{-y\frac{a}{3}}} da_3$$

Is a decreasing function of $s_3$ when $P_3^* = [g^{-3}]_3(b_3) < [f_3^{-3}]_3(r_3)$, that is $[f_3]_3(P_3^*) < r$, $R\big(f_3(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_3(g^{-3}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3\left(g_3^{-3}(b)\right)a_3} > 1. \quad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_3^*(a_3), Q_3^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_3^*(a_3), Q_3^*)$ is more complicated. We again consider perturbations $m_3(t_3 a_3) = p_3(t_3, a_3) - p_3^*(a_3)$ and $W_3(t) = Q_3(t) - Q^*$.Then

$$M_3(t) \equiv \int_0^\infty m_3(t, a_3) da_3 = P_3(t) - P_3^*,$$

$$[m_3]_3 + [m_a]_3 = [p_3]_3 + [p_a]_3 - \left[p^*\right]_3(a_3) \approx -(\mu_3(a_3) + [f_{3*}]_3 + [f_{2*}]_3)m_3$$

$$-[f_{3*}^3]_3 p_3^*(a_3)M - [f_{2*}^3]_3 p_3^*(\mathrm{a})W_3,$$

$$m_3(t,0)p_3(t,0) - p_3^*(0) \approx h_* \int_0^\infty \beta_3(a_3) m_3(t, a_3)\, da + \frac{[h_*^3]_3}{[h_*]_3}\, p_3^*(0) W_3$$

$$\frac{dW_3}{dt} = \frac{dQ_3}{dt} \approx [g_*^3]_3 Q_3^* M_{,3}$$

Where $[f_{3*}]_3$ denoting $[f_3(P^*)]_3$, $[f_{3*}]_3$ denoting $[f_2(P^*)]_3$, $\left[f_{3*}^{'}\right]_3$ denoting $\left[f^{'}{}_3(P^*)\right]_3 \left[f_{2*}^{'}\right]_3$ denoting $[f_2]_3$ $(Q_3^*,$ $h_3^*$ denoting $h_*^{'}$ denoting $h_3^{'}$ $(Q_3^*)$, and $\left[g_*^{'}\right]_3$

Denoting $\left[g^{'}\right]_3(P_3^*)$.

Now, we study the asymptotic behavior of the solutions $m(t, a)$ and $W(t)$ of the system of equations:

$$[m_3]_3 + [m_a]_3 = -(\mu_3(a_3) + [f_{3*}]_3 + [f_{2*}]_3)m_3 - [f'_{3*}]_3[p_3^*](a_3)M_3 - [f'_{2*}]_3 p_3^*(a_3)W_3$$

$$m_3(t,0) = [h_*]_3 \int_0^\infty \beta_3(a_3)m_3(t,a_3)da_3 + \frac{[h_*^3]_3}{[h^*]_3}[p^*]_3(0)\,W_{,3} \qquad (30)$$

$$\frac{dW_3}{dt} = [g_*^3]_3[Q_3^*]M_3$$

Integrating $[m_3]_3(t,a_3)$ with respect to a from 0 to ∞, we have

$$\frac{dM_3(t)}{dt} = \int_0^\infty [mt]_3(t,a_3)da_3$$

$$= -\int_0^\infty [m_a]_3\,(t,a_3)da_3 - \int_0^\infty \mu(a_3)m(t,a_3)da_3$$

$$-\left([f_{3*}]_3 + [f_{2*}]_3 + [f'_{3*}]_3\int_0^\infty p_3^*(a_3)da_3\right)M_3 - [f_{2*}^3]_3\int_0^\infty p_3^*(a_3)da_3 W_3$$

$$= m_3(t,0) - \int_0^\infty \mu_3(a_3)\,m_3(t,a_3)da_3$$

$$-\left([f_{3*}]_3 + [f_{2*}]_3 + [f'_{3*}]_3 P_3^*\right)M_3 - [f'_{2*}]_3 P_3^* W_{,3}$$

From (30), we arrive at the equations

$$\frac{dW_3(t)}{dt} = \int_0^\infty ([h_*]_3\beta_3(a_3) - \mu_3(a_3))m_3(t,a_3)da_3$$

$$-\left([f_{3*}] + [f_{2*}]_3 + [f'_{3*}]_3 P_3^*\right)M_3 + \left(\frac{[h'_*]_3}{[h_*]^3}\,P_3^*(0) - f'_{2*}P_3^*\right)W_3,$$

(33)

$$\frac{dW_3(t)}{dt}\left[g^{'}{}_{*}\right]_3 Q_3^{*} M_{.3}$$

By the mean value theorem,

$$\int_0^{\infty}([h_{*}]_3\beta_3(a_3)-\mu_3(a_3))m_3(t,a_3)da_3=([h_{*}]_3\beta_3(\theta_3)-\mu_3(\theta_3))M_{,3}$$

Where $0<\ \theta_3=\theta_3(t)<\infty$.

And if $a_3^2\,[a_3^2]_3<4\,[a_2]_3<4\,a_2 e$,

$\mathrm{Re}\,\lambda_3\left[\lambda_3\right]_3=\mathrm{Re}\lambda_2\left[\lambda_2\right]_3=-\frac{3}{2}a_3[a2\ ]_3<0.$

Let $\Delta=\max_{2<u<\infty}\left|h_{*}\,[h_{*}]_3\beta_3(a_3)-\mu_3(a_3)\right|$. Then$|n_3(t)|$ 2Δ∀t 0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^{*}(a),\boldsymbol{Q}^{*})$ is asymptotically stable if either $a_3^2<4\,a_2 e$ and $[a_3]>2\Delta$,or $[a_3^2]$ 4$[a_2]$e and $(a_3-\sqrt{(a_3^2)_3-4(a_2)_3 e})>2\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a)=\mu_0+\left(\frac{a}{a+3}\right)$ and $\beta_3(a_3)=ae^{-ma}$ Then $\Delta=[\mu_0]+1$ Simple calculation leads to

$$\int_0^{\infty} \beta_3(a_3)\mathrm{e}^{-\int_0^{\mathrm{a}} \mu_3(\tau_3)+d\tau_3 \mathrm{da}} = \frac{3}{(\mathrm{m}+\mu_0+1)^2}+\frac{2}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_3 + [\mu_0]$ 0.5. Assume$[f_3](x_3)= K_3x_3$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_3(x_3)= [0x_3]/(3 + x_3)$; the functional responses $[[f_2](x) = [0]\, x_3/(1 + x_3)$and $h_3(x_3)$=$(1+\omega_3 x_3)/(1 + x_3)$(See Holling [7] and Rashevsky [13])

Equation (34),now has the form of

$$\frac{m_3 + [\mu_0] + [\mathrm{K}][P_0])}{[m_3]_3 + (\mu_0)_3 + [\mathrm{K}][P_0])_3^3\ [\mathrm{K}]_3[P_0]_3^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(3 + P_0) < b_3$i e $[P_0] < b/(\delta_3 - b_3)$then the equilibrium$(P_0)$ $[P_0]_3(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_3$

$[f_{3*}] = K_3\ \mathrm{P}^*\ [f_{2*}] = \frac{\sigma \mathrm{Q}_3^*}{3+\mathrm{Q}_3^*}$ $[h_*] = \frac{3+\omega_3 \mathrm{Q}_3^*}{3+\mathrm{Q}_3^*}$

$[f_{3*}] = K_3\ [\mathrm{f}_{2*}^3] = \frac{\sigma_3}{(3+\mathrm{Q}_3^*)^2}[\ \mathrm{g}_*^3] = \frac{\sigma_3}{(3+\mathrm{P}_3^*)^2}$ $[\mathrm{h}_*^3] = \frac{\omega_3-3}{(3+\mathrm{Q}_3^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_3}{(m+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma)^2}+\frac{2\omega_3}{([m_3]_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma_3)^3}<1$$

And

$$\frac{1}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^2}+\frac{2\omega_3}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}+(m_3+[\mu_0]_3+\frac{K_3b_3}{\delta_3-b_3}+\sigma_3)\,Q_3^*)^3$$

$$=m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}(1+Q^*)(1+\omega_3 Q_3^*$$

Now

$$\left[a_3^2\right]-4\left[a_2\right]e=\left(\frac{2K_3b_3}{\delta_3-b_3}\right)^2+\frac{4K_3b_3\sigma_3\,Q_3^*}{(\delta_3-b_3)(1+Q_3^*)}+\frac{\sigma_3^2\,Q_3^{*2}}{(1+Q_3^*)^2}$$

$$\frac{4\,b_3\sigma_3\,(\delta_3-b)Q^*}{\delta_3\,(1+Q_3^*)^2}-\frac{4\,(\delta_3-b_3)^2(1-\omega_3)Q_3^*}{\delta_3\,(1+\omega_3Q_3^*)(1+Q_3^*)}P^*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_3\,ae_3^{mu}\;\;[\mu_0]_3-\frac{a_3}{a_3+3}$$

Hence if $[a_3^2>4(a_2)_3e$ and

$$[a_3]=\frac{2K_3b_3}{\delta_3-b_3}+\frac{\sigma_3\,Q_3^*}{1+Q_3^*}>2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_3]_3 = [f_{3*}]_3 + [f_{2*}]_3 + [f_{3*}]P^*_3$$

$$[a_2]_3 = \left[f^3_{2*}\right]P^*_3 - \frac{h^3_*}{h_*}P^*(0)$$

$$e=[g^3_*]_3 Q^*_3$$

$$\eta_3(t) = \left[h_*\right]_3 \beta_3\left(\theta_3(t)\right) - \mu_3\left(\theta(t)\right),$$

The system(33)can be written as the following system:

$$[M(t)]\left[a_3\right]M_3 - [a_2]_3 W_3 - [n(t)]M_3$$

$$w_3(t) = eM_3$$

Write

$$A=\begin{pmatrix} -[a_3]_3 - [a_2]_3 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_3 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_3 \\ W_3 \end{pmatrix}$$

The solution of (32)is

$$X_3(t) = e^{A,t}\,[X]_3(0) + \int_0 e^{At}$$

**This equation is useful in following type of problems . For actual model kindly see various papersThere is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant Λ, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive Λ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the Λ>0 generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible. Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of**

**(e) GW signals from BBHs with component spins χ1,2∈[−0.99,0.99] aligned (e&eb) with the orbital angular momentum, component masses m1,2∈[2,48]M⊙, and (e&eb) total mass Mtotal≤50 M⊙. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins >+0.6 and component masses ≳20 M⊙. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201 STOP Advanced LIGO may be the first experiment to (e) detect gravitational waves. Through superradiance of stellar black holes, it may also be (=) the first experiment to discover the QCD axion with (e&eb) decay constant above the GUT scale. When an axion's Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e) the black hole, forming a "gravitational atom." Through the superradiance process, the number of axions occupying (e) the bound levels grows exponentially, extracting (e) energy and angular momentum from the black hole. Axions transitioning between levels of the gravitational atom and (e&eb) axions annihilating to (e) gravitons can produce (eb) observable gravitational wave signals. The signals are long-lasting, monochromatic, and can be distinguished from (e) ordinary astrophysical sources. Authors estimate up to O (1) transition events at aLIGO for (e) an axion between 10^-11 and (e&eb) 10^-10 eV and up to 10^4 annihilation events for (e) an**

**axion between 10^-13 and 10^-11 eV. In the event of a null search, aLIGO can constrain (e) the axion mass for (e) a range of rapidly spinning black hole formation rates. Axion annihilations are also promising for (e) much lighter masses at future lower-frequency gravitational wave observatories; the rates have (e) large uncertainties, dominated by (e) supermassive black hole spin distributions. Projections for aLIGO are (=) robust against perturbations from (e) the black hole environment and account for (e) updated exclusion on (e&eb) the QCD axion of 6*10^-13 eV < ma < 2*10^-11 eV suggested by (e) stellar black hole spin measurements. Subjects: High Energy Physics - Phenomenology (hep-ph); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 91, 084011 (2015) DOI: 10.1103/PhysRevD.91.084011Cite as:arXiv:1411.2263 [hep-ph] (or arXiv:1411.2263v3 [hep-ph] for this version) Discovering the QCD Axion with Black Holes and Gravitational Waves Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang STOP Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by**

**(e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave signal. The signal sweeps upwards in frequency from 35 to 250 Hz with (e&eb) a peak gravitational-wave strain of (e) 1.0×10−21. It matches the waveform predicted by general relativity for (e) the inspiral and merger of a pair of black holes and the ringdown of the resulting (eb) single black hole. The signal was observed with a matched-filter signal-to-noise ratio of (e) 24 and a false alarm rate estimated to be (=) less**

**than 1 event per 203 000 years, equivalent to (=) a significance greater than 5.1σ. The source lies at (eb) a luminosity distance of 410+160−180 Mpc corresponding to (e&eb) a redshift z=0.09+0.03−0.04. In the source frame, the initial black hole masses are (=) 36+5−4M⊙ and 29+4−4M⊙, and the final black hole mass is (=) 62+4−4M⊙, with (e&eb) 3.0+0.5−0.5M⊙c2 radiated in (eb) gravitational waves. All uncertainties define (eb) 90% credible intervals. These observations demonstrate (eb) the existence of binary stellar-mass black hole systems. Observation of Gravitational Waves from a Binary Black Hole Merger B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 061102 – Published 11 February 2016 This is the first direct detection of gravitational waves and (e&eb) the first observation of a binary black hole merger. Received 21 January 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.061102 STOP This is the second in a series of three papers on (e&eb) systematic analysis of rank 1 Coulomb branch geometries of (e) four dimensional N=2 SCFTs. In the first paper authors developed a strategy for (e) classifying physical rank-1 CB geometries of N=2 SCFTs. They show how to carry out this strategy computationally to construct (eb) the Seiberg-Witten curves and (e&eb) one-forms for all the rank-1 SCFTs. Explicit expressions are given for all cases, with the exception of the Nf=4 SU (2) gauge theory and (e&eb) the En SCFTs which were previously constructed. Classification includes all known rank-1 theories plus (+) a new one with (e&eb) an abelian flavor group, plus (+) nine additional theories whose existence is more speculative. Four of those, reported in first paper, depend on (=) the assumption of new frozen rank-1 SCFTs. Here authors also show that the assumption of the existence of certain rank-0 N=2 SCFTs leads to (eb) five additional consistent rank-1 CB geometries. Subjects: High Energy Physics - Theory (hep-th) Cite as:arXiv: 1601.00011 [hep-th] (or arXiv: 1601.00011v2 [hep-th] for this version Geometric constraints on the space of N=2 SCFTs II: Construction of special Kähler geometries and RG flows Philip C. Argyres, Matteo Lotito, Yongchao Lü, Mario Martone (Submitted on 31 Dec 2015 (v1), last revised 6 Jan 2016 (this version, v2)) STOP Authors**

**propose a model of Dark Energy in which the field currently dominating (e) the energy density of the universe is (=) an "axion field" linearly coupled to (e&eb) the Pontryagin density, tr (FΛF), (i.e., the exterior derivative of (e) the Chern-Simons form) of (e) a massive gauge field. Authors assume that the axion has (e) self-interactions corresponding to (e&eb) a non-trivial (exponential) potential. Authors argue that a non-vanishing magnetic helicity of (e) the gauge field triggers (eb) slow-rolling of the axion at field values far below (eb) the Planck scale. Proposal leads to (eb) a "Tracking Dark Energy Scenario" in which the contribution of the axion energy density to (e) the total energy density is (=) constant (and small) during (e&eb) the early radiation phase, until (e) a secular growth term proportional to (e) the Pontryagin density of (e) the gauge field becomes (=) dominant. The initially small contribution of (e) the axion field to (e) the total energy density is related to the observed small baryon-to-entropy ratio. Some speculations concerning the nature of the gauge field are offered.**

**Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Cite as:arXiv: 1601.00057 [hep-th] (or arXiv:1601.00057v2 [hep-th] for this version) Tracking Dark Energy from Axion-Gauge Field Couplings Stephon Alexander, Robert Brandenberger, Juerg Froehlich (Submitted on 1 Jan 2016 (v1), last revised 23 Jan 2016 (this version, v2)) STOP Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=)**

**Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version) Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP In this work authors obtain bounds on (e&eb) the topological Abelian string-vortex and on (e&eb) the string-cigar, by using (e) a new measure of configurational complexity, known as (=) configurational entropy. In this way, the information-theoretical measure of (e) six-dimensional braneworld scenarios is capable to (e) probe situations where (e) the parameters responsible for (e) the brane thickness are arbitrary. The so-called configurational entropy (CE) selects (e&eb) the best value of the parameter in (eb) the model. This is accomplished by (e) minimizing the CE, namely, by (e) selecting the most appropriate parameters in the model that correspond to (e&eb) the most organized system, based upon (e) the Shannon information theory. This information-theoretical measure of complexity provides (eb) a complementary perspective to (e) situations where (e) strictly energy-based arguments are (=) inconclusive. Authors show that the higher the energy the higher the CE, what shows (eb) an important correlation between the energy of the localized field configuration and (e&eb) its associated entropic measure. Subjects: High Energy Physics - Theory (hep-th) DOI:10.1016/j.physletb.2016.02.038 Cite as: arXiv: 1601.00076 [hep-th] (or arXiv: 1601.00076v2 [hep-th] for this version) Bounds on topological Abelian string-vortex and string-cigar from information-entropic measure R. A. C. Correa, D. M. Dantas, C. A. S. Almeida, Roldao da Rocha (Submitted on 1 Jan**

**2016 (v1), last revised 17 Feb 2016 (this version, v2)) STOP Sum of the Holevo quantity (that bounds the capacity of quantum channels to transmit (e&eb) classical information about an observable) and (e&eb) the quantum discord (a measure of the quantumness of (e) correlations of that observable) yields (eb) an observable-independent total given by (e) the quantum mutual information This split naturally delineates ,dilates and disseminates information about (e&eb) quantum systems accessible to (e) observers -- information that is redundantly transmitted by (e&eb) the environment -- while showing that it is (=) maximized for (e) the quasi-classical pointer observable. Other observables are accessible only via (e&eb) correlations with the pointer observable .Further authors prove an anti-symmetry property relating accessible information and (e&eb) discord. It shows that information becomes (=) objective -- accessible to (e) many observers -- only as quantum information is relegated to (e) correlations with (e&eb) the global environment, (eb) and, therefore, locally inaccessible. The resulting complementarity explains (eb) why, in a quantum Universe, we perceive objective classical reality while (e) flagrantly quantum superpositions are out of reach. Subjects: Quantum Physics (quant-ph) Journal reference: Scientific Reports 3, 1729 (2013) Report number: LAUR-13-26386 Cite as:arXiv: 1303.4659 [quant-ph] (or arXiv: 1303.4659v2 [quant-ph] for this version) Complementarity of quantum discord and classically accessible information Michael Zwolak, Wojciech H. Zurek STOP Niels Bohr proposed that the outcome of the measurement becomes (=) objective and real, and, hence, classical, when (e) its results can be communicated by (e) classical means. In this work authors revisit Bohr's postulate using (e) modern tools from (e) the quantum information theory. They find a full confirmation of Bohr's idea: if a measurement device is in a nonclassical state, (eb) the measurement results cannot be communicated perfectly by classical means. In this case some part of information in the measurement apparatus is lost (e) in the process of communication: the amount of this lost information turns out to be (=) the quantum discord. The information loss occurs even when (e) the apparatus is not entangled**

**with the system of interest. The tools presented in this work allow (eb) to generalize Bohr's postulate: authors show that for pure system-apparatus states (eb) quantum communication does not provide any advantage when (e) measurement results are communicated to (e) more than one recipient. They further demonstrate the superiority of quantum communication to (e&eb) two recipients on (e&eb) a mixed system-apparatus state and show that (eb) this effect is fundamentally different from (e) quantum state cloning. Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. Lett. 111, 040401 (2013) DOI:10.1103/PhysRevLett.111.040401 Report number: LA-UR-12-25351 Cite as: arXiv: 1304.2914 [quant-ph] (or arXiv: 1304.2914v2 [quant-ph] for this version) Quantum discord cannot be shared Alexander Streltsov, Wojciech H. Zurek STOP Wojciech H. Zurek, Don N. Page investigate thermodynamic equilibrium of (e) a self-gravitating perfect fluid in (eb) a spherically symmetric system containing (e) a black hole of mass M by means of (e) the Tolman-Oppenheimer-Volkoff (TOV) equation. At r >> 2M its solutions describe (eb) a black-body radiation atmosphere with (e&eb) the Hawking temperature T_BH~1/(8 \pi M) that is increasingly blueshifted as r approaches (eb) 2M. However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments:**

**Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by (e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a firewall-like "Planck cocoon" with entropy close to (e) black hole entropy**

## Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

# Module Four

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equailibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-445) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Some of the concepts to which model is applicable is as follows:**

**Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by (e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave signal. The signal sweeps upwards in frequency from 35 to 250 Hz with (e&eb) a peak gravitational-wave strain of (e) 1.0×10−21. It matches the waveform predicted by general relativity for (e) the inspiral and merger of a pair of black holes and the ringdown of the resulting (eb) single black hole. The signal was observed with a matched-filter signal-to-noise ratio of (e) 24 and a false alarm rate estimated to be (=) less than 1 event per 203 000 years, equivalent to (=) a significance greater than 5.1σ. The source lies at (eb) a luminosity distance of 410+160−180 Mpc corresponding to (e&eb) a redshift z=0.09+0.03−0.04. In the source frame, the initial black hole masses are (=) 36+5−4M⊙ and 29+4−4M⊙, and the final black hole mass is (=) 62+4−4M⊙, with (e&eb) 3.0+0.5−0.5M⊙c2 radiated in (eb) gravitational waves. All uncertainties define (eb) 90% credible intervals. These observations demonstrate (eb) the existence of binary stellar-mass black hole systems. Observation of Gravitational Waves from a Binary Black Hole Merger B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 061102 – Published 11 February 2016 This is the first direct detection of gravitational waves and (e&eb) the first observation of a binary black hole**

**merger. Received 21 January 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.061102 STOP This is the second in a series of three papers on (e&eb) systematic analysis of rank 1 Coulomb branch geometries of (e) four dimensional N=2 SCFTs. In the first paper authors developed a strategy for (e) classifying physical rank-1 CB geometries of N=2 SCFTs. They show how to carry out this strategy computationally to construct (eb) the Seiberg-Witten curves and (e&eb) one-forms for all the rank-1 SCFTs. Explicit expressions are given for all cases, with the exception of the Nf=4 SU (2) gauge theory and (e&eb) the En SCFTs which were previously constructed. Classification includes all known rank-1 theories plus (+) a new one with (e&eb) an abelian flavor group, plus (+) nine additional theories whose existence is more speculative. Four of those, reported in first paper, depend on (=) the assumption of new frozen rank-1 SCFTs. Here authors also show that the assumption of the existence of certain rank-0 N=2 SCFTs leads to (eb) five additional consistent rank-1 CB geometries. Subjects: High Energy Physics - Theory (hep-th) Cite as:arXiv: 1601.00011 [hep-th] (or arXiv: 1601.00011v2 [hep-th] for this version Geometric constraints on the space of N=2 SCFTs II: Construction of special Kähler geometries and RG flows Philip C. Argyres, Matteo Lotito, Yongchao Lü, Mario Martone (Submitted on 31 Dec 2015 (v1), last revised 6 Jan 2016 (this version, v2)) STOP Authors propose a model of Dark Energy in which the field currently dominating (e) the energy density of the universe is (=) an "axion field" linearly coupled to (e&eb) the Pontryagin density, tr (F∧F), (i.e., the exterior derivative of (e) the Chern-Simons form) of (e) a massive gauge field. Authors assume that the axion has (e) self-interactions corresponding to (e&eb) a non-trivial (exponential) potential. Authors argue that a non-vanishing magnetic helicity of (e) the gauge field triggers (eb) slow-rolling of the axion at field values far below (eb) the Planck scale. Proposal leads to (eb) a "Tracking Dark Energy Scenario" in which the contribution of the axion energy density to (e) the total energy density is (=) constant (and small) during (e&eb) the early radiation phase, until (e) a secular growth term proportional to (e) the Pontryagin density of (e) the gauge field becomes (=) dominant. The initially small contribution of (e) the axion field to (e) the total energy density is related to the observed small baryon-to-entropy ratio. Some speculations concerning the nature of the gauge field are offered.**
**Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Cite as:arXiv: 1601.00057 [hep-th] (or arXiv:1601.00057v2 [hep-th] for this version) Tracking Dark Energy from Axion-Gauge Field Couplings Stephon Alexander, Robert Brandenberger, Juerg Froehlich (Submitted on 1 Jan 2016 (v1), last revised 23 Jan 2016 (this version, v2)) STOP Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version) Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP In this work authors obtain bounds on (e&eb) the topological Abelian string-vortex and on (e&eb) the string-cigar, by using (e) a new measure of configurational complexity, known as (=) configurational entropy. In this way, the information-theoretical measure of (e) six-dimensional braneworld scenarios is capable to (e) probe situations where (e) the parameters responsible for (e) the brane thickness are arbitrary. The so-called configurational entropy (CE) selects (e&eb) the best value of the parameter in (eb) the model. This is accomplished by (e) minimizing the CE, namely, by (e) selecting the most appropriate parameters in the model that correspond to (e&eb) the most organized system, based upon (e) the Shannon information theory. This information-theoretical measure of complexity provides (eb) a complementary perspective to (e) situations where (e) strictly energy-based arguments are (=) inconclusive. Authors show that the higher the energy the higher the CE, what shows (eb) an important correlation between the energy of the localized field configuration and (e&eb) its**

**associated entropic measure. Subjects: High Energy Physics - Theory (hep-th) DOI:10.1016/j.physletb.2016.02.038 Cite as: arXiv: 1601.00076 [hep-th] (or arXiv: 1601.00076v2 [hep-th] for this version) Bounds on topological Abelian string-vortex and string-cigar from information-entropic measure R. A. C. Correa, D. M. Dantas, C. A. S. Almeida, Roldao da Rocha (Submitted on 1 Jan 2016 (v1), last revised 17 Feb 2016 (this version, v2)) STOP Sum of the Holevo quantity (that bounds the capacity of quantum channels to transmit (e&eb) classical information about an observable) and (e&eb) the quantum discord (a measure of the quantumness of (e) correlations of that observable) yields (eb) an observable-independent total given by (e) the quantum mutual information This split naturally delineates ,dilates and disseminates information about (e&eb) quantum systems accessible to (e) observers -- information that is redundantly transmitted by (e&eb) the environment -- while showing that it is (=) maximized for (e) the quasi-classical pointer observable. Other observables are accessible only via (e&eb) correlations with the pointer observable .Further authors prove an anti-symmetry property relating accessible information and (e&eb) discord. It shows that information becomes (=) objective -- accessible to (e) many observers -- only as quantum information is relegated to (e) correlations with (e&eb) the global environment, (eb) and, therefore, locally inaccessible. The resulting complementarity explains (eb) why, in a quantum Universe, we perceive objective classical reality while (e) flagrantly quantum superpositions are out of reach. Subjects: Quantum Physics (quant-ph) Journal reference: Scientific Reports 3, 1729 (2013) Report number: LAUR-13-26386 Cite as:arXiv: 1303.4659 [quant-ph] (or arXiv: 1303.4659v2 [quant-ph] for this version) Complementarity of quantum discord and classically accessible information Michael Zwolak, Wojciech H. Zurek STOP Niels Bohr proposed that the outcome of the measurement becomes (=) objective and real, and, hence, classical, when (e) its results can be communicated by (e) classical means. In this work authors revisit Bohr's postulate using (e) modern tools from (e) the quantum information theory. They find a full confirmation of Bohr's idea: if a measurement device is in a nonclassical state, (eb) the measurement results cannot be communicated perfectly by classical means. In this case some part of information in the measurement apparatus is lost (e) in the process of communication: the amount of this lost information turns out to be (=) the quantum discord. The information loss occurs even when (e) the apparatus is not entangled with the system of interest. The tools presented in this work allow (eb) to generalize Bohr's postulate: authors show that for pure system-apparatus states (eb) quantum communication does not provide any advantage when (e) measurement results are communicated to (e) more than one recipient. They further demonstrate the superiority of quantum communication to (e&eb) two recipients on (e&eb) a mixed system-apparatus state and show that (eb) this effect is fundamentally different from (e) quantum state cloning. Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. Lett. 111, 040401 (2013) DOI:10.1103/PhysRevLett.111.040401 Report number: LA-UR-12-25351 Cite as: arXiv: 1304.2914 [quant-ph] (or arXiv: 1304.2914v2 [quant-ph] for this version) Quantum discord cannot be shared Alexander Streltsov, Wojciech H. Zurek STOP Wojciech H. Zurek, Don N. Page investigate thermodynamic equilibrium of (e) a self-gravitating perfect fluid in (eb) a spherically symmetric system containing (e) a black hole of mass M by means of (e) the Tolman-Oppenheimer-Volkoff (TOV) equation. At r >> 2M its solutions describe (eb) a black-body radiation atmosphere with (e&eb) the Hawking temperature T_BH~1/(8 \pi M) that is increasingly blueshifted as r approaches (eb) 2M. However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments: Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by (e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a firewall-like "Planck cocoon" with entropy close to (e) black hole entropy. Authors reproduce paper with an updated title and abstract. For a later study with similar results see G. 't Hooft, gr-qc/9706058 Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Physical Review D, Volume 29, Number 4, pp. 628 - 631 (1984) DOI:10.1103/PhysRevD.29.628 Cite as:arXiv:1511.07051 [gr-qc] (or arXiv:1511.07051v1 [gr-qc] for this version) Firewalls, black-hole thermodynamics, and singular solutions of the Tolman-Oppenheimer-Volkoff equation Wojciech H. Zurek, Don N. Page (Submitted on 22 Nov 2015) STOP Frederic Magniez, Ashwin Nayak consider the problem of testing the commutativity of (e) a black-box group specified by its k generators. The complexity (in terms of k) of (e) this problem was first considered by Pak, who gave (eb) a**

**randomized algorithm involving O(k) group operations. They construct a quite optimal quantum algorithm for this problem whose complexity is in O~(k2/3)O~(k2/3). The algorithm uses and highlights the power of the quantization method of Szegedy. For the lower bound of Ω(k2/3)Ω(k2/3), authors give a reduction from (e) a special case of Element Distinctness to our problem. Along the way, they prove the optimality of the algorithm of (e) Pak for the randomized model. Algorithmica July 2007, Volume 48, Issue 3, pp 221-232 First online: 25 June 2007 Quantum Complexity of Testing Group Commutativity Frederic Magniez, Ashwin Nayak STOP Pawel Wocjan, Jon Yard analyze relationships between quantum computation and (e&eb) a family of generalizations of the Jones polynomial. Extending recent work by Aharonov et al., authors give efficient quantum circuits for (e) implementing the unitary Jones-Wenzl representations of the braid group. We use these to provide new quantum algorithms for (e) approximately evaluating a family of specializations of the HOMFLYPT two-variable polynomial of trace closures of braids. They also give (eb) algorithms for approximating the Jones polynomial of a general class of closures of (e) braids at roots of unity. Next authors provide (eb) a self-contained proof of a result of (e) Freedman et al. that any quantum computation can be replaced by (e&eb) an additive approximation of the Jones polynomial, evaluated at (eb) almost any primitive root of unity. Proof encodes two-qubit unitaries into (e&eb) the rectangular representation of (e) the eight-strand braid group. Authors then give QCMA-complete and (e&eb) PSPACE-complete problems which are based on braids. They conclude with direct proofs that evaluating the Jones polynomial of the plat closure at most primitive roots of unity is a #P-hard problem, while learning its most significant bit is PP-hard, circumventing (e) the usual route through the Tutte polynomial and graph coloring. Subjects: Quantum Physics (quant-ph); Geometric Topology (math.GT) Cite as: arXiv:quant-ph/0603069 (or arXiv:quant-ph/0603069v3 for this version) The Jones polynomial: quantum algorithms and applications in quantum complexity theory Pawel Wocjan, Jon Yard STOP Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan address the task of switching off the Hamiltonian of (e) a system by removing (e) all internal and system-environment couplings. Authors propose dynamical decoupling schemes, which use (e) only bounded-strength controls, for quantum many-body systems with local system Hamiltonians and local environmental couplings. To do so, they introduce (eb) the combinatorial concept of (e) balanced-cycle orthogonal arrays (BOAs) and show how to construct them from (e) classical error-correcting codes. The derived decoupling schemes may be useful as (=) a primitive for more complex schemes, e.g., for Hamiltonian simulation. For the case of n qubits and a 2-local Hamiltonian, the length of the resulting decoupling (e&eb) scheme scales as O (nlogn), improving over (e&eb) the previously best-known schemes that scaled quadratically with (e&eb) n. More generally, using balanced-cycle orthogonal arrays constructed from families of BCH codes, we show that bounded-strength decoupling for any ℓ-local Hamiltonian, where ℓ≥2, can be achieved using decoupling schemes of length at most O(nℓ−1logn). Subjects:Quantum Physics (quant-ph); Emerging Technologies (cs.ET) Journal reference: IEEE Trans. Inf. Theory, 62(5): 2881-2894, 2016 DOI: 10.1109/TIT.2016.2535183 Report number: MIT-CTP /4705 Cite as: arXiv: 1509.00408 [quant-ph] (or arXiv: 1509.00408v2 [quant-ph] for this version) Improved bounded-strength decoupling schemes for local Hamiltonians Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan STOP Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola propose dynamical control schemes for Hamiltonian simulation in many-body quantum systems that avoid instantaneous control operations and rely solely on (e) realistic bounded-strength control Hamiltonians. Each simulation protocol consists of (e) periodic repetitions of a basic control block, constructed as (=) a suitable modification of an "Eulerian decoupling cycle," that would otherwise implement (eb) a trivial (zero) target Hamiltonian. For an open quantum system coupled to (e) an uncontrollable environment, approach may be employed to (e) engineer an effective evolution that simulates (e&eb) a target Hamiltonian on the system, while suppressing (e) unwanted decoherence to the leading order. Authors present illustrative applications to (e) both closed- and open-system simulation settings, with (e&eb) emphasis on simulation of non-local (two-body) Hamiltonians using (e) only local (one-body) controls. In particular, authors provide (eb) simulation schemes applicable to (e&eb) Heisenberg-coupled spin chains exposed to (e) general linear decoherence, and show how to simulate (e&eb) Kitaev's honeycomb lattice Hamiltonian starting from Ising-coupled qubits, as (=) potentially relevant to (e&eb) the dynamical generation of a topologically protected (e&eb) quantum memory. Additional implications for quantum information processing are discussed. Subjects: Quantum Physics (quant-ph) Journal reference: New J. Phys. 16, 045021 (2014) DOI: 10.1088/1367-2630/16/4/045021 Cite as:arXiv: 1310.4153 [quant-ph] (or arXiv: 1310.4153v1 [quant-ph] for this version) Hamiltonian quantum simulation with bounded-strength controls Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola STOP Felix Fontein, Pawel Wocjan**

**present a quantum algorithm for computing the period lattice of (e) infrastructures of fixed dimension. The algorithm applies to (e&eb) infrastructures that satisfy certain conditions. The latter are always fulfilled for (e) infrastructures obtained from (e) global fields, i.e., algebraic number fields and function fields with (e&eb) finite constant fields. The first main contributions is an exponentially better method for (e) sampling approximations of vectors of the dual lattice of (e) the period lattice than the methods outlined in the works of (e) Hallgren and Schmidt and Vollmer. This new method improves (eb) the success probability by (e) a factor of at least 2^ {n^2-1} where (e) n is the dimension. The second main contribution is a rigorous and complete proof that the running time of the algorithm is (=) polynomial in the logarithm of (e) the determinant of the period lattice and (e&eb) exponential in n. The third contribution is the determination of (e) an explicit lower bound on (e&eb) the success probability of our algorithm which greatly improves on (e&eb) the bounds given in the above works. The exponential scaling seems inevitable because (e) the best currently known methods for (e) carrying out fundamental arithmetic operations in infrastructures obtained from (e) algebraic number fields take exponential time. In contrast, the problem of computing (e&eb) the period lattice of infrastructures arising from function fields can be solved without (e) the exponential dependence on (e&eb) the dimension n since this problem reduces efficiently to the abelian hidden subgroup problem. This is also true for (e) other important computational problems in (eb) algebraic geometry. The running time of the best classical algorithms for (e) infrastructures arising from (e) global fields increases subexponentially with (e&eb) the determinant of the period lattice. Subjects: Quantum Physics (quant-ph) Cite as: arXiv: 1111.1348 [quant-ph] (or arXiv: 1111.1348v2 [quant-ph] for this version) Quantum Algorithm for Computing the Period Lattice of an Infrastructure Felix Fontein, Pawel Wocjan STOP Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan advocate a new approach of (e) addressing hidden structure problems and (e&eb) finding efficient quantum algorithms. Authors introduce and investigate the Hidden Symmetry Subgroup Problem (HSSP), which is a generalization of the well-studied Hidden Subgroup Problem (HSP). Given a group acting on (e&eb) a set and an oracle whose level sets define (eb) a partition of the set; the task is to recover (e) the subgroup of symmetries of (e) this partition inside the group. The HSSP provides (eb) a unifying framework that, besides the HSP, encompasses (e) a wide range of algebraic oracle problems, including (e) quadratic hidden polynomial problems. While the HSSP can have (e) provably exponential quantum query complexity, authors obtain efficient quantum algorithms for various interesting cases. To achieve this, authors present a general method for (e) reducing the HSSP to the HSP, which works efficiently in several cases related to (e&eb) symmetries of polynomials. The HSSP therefore connects in a rather surprising way certain hidden polynomial problems with (e&eb) the HSP. Using this connection, authors obtain the first efficient quantum algorithm for (e) the hidden polynomial problem for (e) multivariate quadratic polynomials over fields of (e) constant characteristic. Authors also apply the new methods to (e&eb) polynomial function graph problems and (e&eb) present an efficient quantum procedure for (e) constant degree multivariate polynomials over (e&eb) any field. This result improves in several ways the currently known algorithms. Subjects: Quantum Physics (quant-ph) Journal reference: SIAM Journal on Computing 42:5 (2013) DOI: 10.1137/120864416 Cite as: arXiv: 1107.2189 [quant-ph] (or arXiv: 1107.2189v2 [quant-ph] for this version) Hidden Symmetry Subgroup Problems Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan STOP Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan present an efficient general method for realizing a quantum walk operator corresponding to (e&eb) an arbitrary sparse classical random walk. Approach is based on Grover and Rudolph's method for preparing coherent versions of (e) efficiently integrable probability distributions. This method is intended for (e) use in quantum walk algorithms with polynomial speedups, whose complexity is usually measured in terms of (e&eb) how many times we have to apply a step of a quantum walk, compared to (e&eb) the number of necessary classical Markov chain steps. Authors consider a finer notion of complexity including (e) the number of elementary gates it takes to implement each step of the quantum walk with (e&eb) some desired accuracy. The difference in complexity for (e) various implementation approaches is that method scales linearly in the sparsity parameter and (e&eb) poly-logarithmically with (e&eb) the inverse of the desired precision. Best previously known general methods either scale quadratically in (eb) the sparsity parameter, or polynomially in (eb) the inverse precision. Approach is especially relevant for (e) implementing quantum walks corresponding to (e&eb) classical random walks like those used in (eb) the classical algorithms for approximating permanents and sampling from binary contingency tables. In those algorithms, the sparsity parameter grows with (e&eb) the problem size, while (e) maintaining high precision is required. Subjects: Quantum Physics (quant-ph) Journal reference: Quantum Information and Computation, Vol.10,**

**No.5&6, pp0420-0434 (2010) Cite as: arXiv: 0903.3465 [quant-ph] (or arXiv: 0903.3465v3 [quant-ph] for this version) Efficient Circuits for Quantum Walks Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan STOP Preparing the ground state of a system of interacting classical particles is (=) an NP-hard problem. Thus, there is in general no better algorithm to (e) solve this problem than exhaustively going through all N configurations of (e) the system to determine (eb) the one with lowest energy, requiring (e) a running time proportional to N. A quantum computer, if it could be built, could solve (e) this problem in time sqrt (N). Here, authors present a powerful extension of this result to (e) the case of interacting quantum particles, demonstrating (eb)that a quantum computer can prepare the ground state of (e) a quantum system as efficiently as it does for (e) classical systems. Subjects: Quantum Physics (quant-ph) Journal reference: Phys.Rev.Lett.102:130503, 2009 DOI:10.1103/PhysRevLett.102.130503 Cite as: arXiv: 0809.2705 [quant-ph] (or arXiv: 0809.2705v1 [quant-ph] for this version) Preparing ground states of quantum many-body systems on a quantum computer David Poulin, Pawel Wocjan STOP Daniel Nagaj, Pawel Wocjan construct a simple translationally invariant, nearest-neighbor Hamiltonian on (e&eb) a chain of 10-dimensional qudits that makes it possible to (e) realize universal quantum computing without (e) any external control during (e&eb) the computational process. Authors only require the ability to prepare an initial computational basis state which encodes (e) both the quantum circuit and its input. The computational process is (=) then carried out by (e) the autonomous Hamiltonian time evolution. After a time polynomially long in the size of (e) the quantum circuit has passed, the result of (e) the computation is obtained with (e&eb) high probability by (e) measuring a few qudits in the computational basis. This result also implies (eb) that there cannot (e) exist efficient classical simulation methods for (e) generic translationally invariant nearest-neighbor Hamiltonians on (e&eb) qudit chains, unless (e) quantum computers can be efficiently simulated by (e&eb) classical computers (or, put in complexity theoretic terms, unless BPP=BQP). Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. A 78, 032311 (2008) DOI: 10.1103/PhysRevA.78.032311 Cite as: arXiv: 0802.0886 [quant-ph] (or arXiv: 0802.0886v4 [quant-ph] for this version) Hamiltonian Quantum Cellular Automata in 1D Daniel Nagaj, Pawel Wocjan**

## 55. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_4}{dt} = -b_4 Q_4 + g_4(g_4)Q_4$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_4(0) = 0, \quad g_4(P_4) > 0, \ P_4 \ \square 0, \qquad \lim_{P_4 \to \infty} g_4(P_4) = (K_4)_4 \ \square\infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_4 + (\rho_a)_4 + D_4(a_4, P_4, Q_4)\rho_4 + 0, \qquad 0 < a_4 ?\infty,$$

$$\rho_4(t,0) = \int_0^{\infty} n_4(a_4, P_4, Q_4)\rho_4(t_4, a_4)da_4$$

Where $D_4$ is the mortality function and $n_4$ is the fertility function.

Put $_4$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_4}{dt}\right]_4 = x_4\left[\frac{dQ_2}{dt}\right]_4 - x_2\left[\frac{dQ_4}{dt}\right]_4$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_4 - x_4\left[\frac{dQ_2}{dt}\right]_4$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_4 - x_6\left[\frac{dQ_3}{dt}\right]_4$$

Where $x_4, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_4$ and $g_4(P_4)Q_4$

Now when you add the term $(b_4Q_4)$ with $g_4(P_4)(Q_4)$ we are infect see is that $(b_4Q_4)$ is credited incrementally part of $g_4(P_4)(Q_4)$ or instant nearly, when subtracted from $\frac{dQ_4}{dt}$ , it means that $g_4(P_4)(Q_4)$ is slowly eaten by $\frac{dQ_4}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_4 = \mu_4(a_4) + (f_4)_4(P_4)_4 + (f_2)_4(Q)_4 \quad (2)$$

$$n_4 = \beta_4(a_4)h_4(Q_4), \quad (3)$$

Where $\mu_4$the age- dependent debit rate whit is $\mu_4(a_4)$ 0 $\forall a > 0$; $(f_4)_4$ is the density- dependent death rate with $f_4(0) = 0$, $(f_4^4)_4(P_4)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_4)_4(P_4) = (k_4)_4$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_4(0) = 0, (f_2^4)_4(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_4(Q_4) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 4, h^4(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_4 + (\rho_a)_4 = -\mu_4(a_4) + (f_4)_4P_4) + (f_2)_4(Q_4))(\rho_4)_4 \quad (4)$$

$$\rho_4(t_4, 0) = h_4(Q_4)\int_0^\infty \beta_4(a_4)\rho_4(t, a_4)da, \quad (5)$$

$$\frac{dQ_4}{dt} = -b_4Q_4 + g_4(P_4)Q_4 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_4$ and $Q_4$ of (4-6) satisfy

$$0 \ \rho_4(t, a_4) \ \rho_4^*(t, a_4),$$

$$0 \ Q_4(t) \ Q_4^*(t),$$

Where $\rho_4^*(t, a_4)$ and $Q_4^*(t)$, are solutions of

$$(\rho_t^*)_4 + (\rho_a^*)_4 = -(\mu_4(a_4) + (f_4)_4(P_4^*) + (f_2)_4Q_4^*))\,\rho_4^*, \quad (7)$$

$$\rho_4^*(t,0) = \int_0^\infty \beta_4(a_4)\,\rho_4^*(t,a_4)da_4 \quad (8)$$

$$\frac{dQ_4^*}{dt} = -b_4\,Q_4^* + g_4(P_4^*)Q_4^*, \quad (9)$$

With $P_4^*(t) = \int_0^\infty \rho_4^*(t,a_4)da_4$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_4^*(t)}{dt} = P_4^*(t)\big(r_4 - (f_4)(P_4^*) - (f_2)_4(Q^*)\big), \quad (10)$$

$$\frac{dQ_4^*(t)}{dt} = Q_4^*(t)(-b_4 + g_4(P_4^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_4(a)e^{-\int_0^a \mu_4(\tau)dt}e^{-ra}\,da_4 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_4^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(t)dt}\,da_4 > 1, \quad (H1)$$

Which guarantees $r_4 > 0$.

*Remark.* If $\beta_4(a) = [\beta_0]_4 a_4 e^{-\alpha_4 a_4}, \mu_4(a_4) \equiv \mu_4$ then the solution of (12) is

$$r_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_4]_4(P_4) = 0, [f_2](Q_4) = \lambda_4 Q_4\ h_4(Q_4) = 1, \qquad and\ \ g_4(P_4) = c_4 P_4$$

$$[\rho_4]_4 + [\rho_a]_4 = -(\mu_4 + \lambda_4 Q_4)\rho_4$$

$$\rho_4(t,0) = \int_0^\infty [\beta_0]_4 a_4 e^{-\alpha_4 a_4} \rho_4(t,a_4)\, da_4$$

$$\frac{dQ_4}{dt} = -b_4 Q_4 + c_4 P_4 Q_4$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_4}{dt} = \omega_4 P_4 - \lambda_4 P_4 Q_4$$

$$\frac{dQ_4}{dt} = b_4 Q_4 - c_4 P_4 Q_4$$

Where $\omega_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$

## 56. THE EQUILIBRIUM $(\rho_0)_4(a_4), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_4(a_4), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_4 + (\rho_a)_4 = -\big(\mu_4(a_4) + (f_4)_4(P_4)\big)\rho,$$

$$\rho_4(t,0) = \int_0^\infty \beta_4(a_4)\rho_4(t,a_4)da_4 \qquad (13)$$

The solution of (13) has the form of

$$\rho_4(t,a_4) = Ce_4^{r(t-a_4)} - \int_0^a (\tau)d\tau - \int_0^4 (f_4)_4\big(P_4(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_4(t) = C_4 e^{r_4 t - \int_0^t (f_4)_4(P_4(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_4^{-4}(r)$ if $k_4 > r$. This is equivalent to

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{(k_4)_4 a_4} da_4 < 1.$$

In fact, if we define

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau)d\tau} e^{-s_4 a_4}\, da_4$$

$R(s)$ is a decreasing function. $R_4(r_4) = 1$ And $R_4(k_4)_4 < 1$ imply $r_4 < (k_4)_4$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_4(a_4), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{k_4 a_4} da_4$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_4 \equiv \int_0^\infty (\rho_0)_{44}(a_4)da_4$ *satisfies*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{-(f_4)_4(P_0)a_4} da$$

$$= 1 \qquad (11)$$

Next, we study the stability of $(\rho_0)_4(a_4), 0)$

Let $u_4(t, a_4) = \rho(t, a_4) - \rho_0(a_4)$ Then

$$N_4(t) \equiv \int_0^\infty u_4(t, a_4)da_4 = \int_0^\infty \rho_4(t, a_4)da_4 - \int_0^\infty (\rho_0)_4(a_4)da = P_4(t) - P_0)_4$$

And

$$(u_4)_4 + (u_a)_4 = (\rho_t)_4 + (\rho_a)_4 - (\rho_0^4)_4(a_4)$$

$$= -\mu_4(a_4) + (f_4)_4(P_4) + (f_2)_4(Q_4))\,(u_4 + (\rho_0)_4) + \mu_4(a_4)(\rho_0)_4 + (f_4)_4(P_0)_4(\rho_0)_4$$

$$= -\mu_4(a_4)u_4 - (f_4)_4(P_4) + (f_2)_4(Q_4))\,u_4 - (\rho_0)_4\,(f_4)_4(P_4) - (f_4)_4(P_0)_4$$

$$+ (f_2)_4(Q_4)), \qquad (15)$$

Where $(\rho_0^4)_4(a_4)$ denotes $d(\rho_0)_4(a_4)/da_4$

Expansion $(f_4)_4(P_4)$ and $(f_2)_4(Q_4)$ at $(P_0)_4$ and 0 yield respectively

$$(f_4)_4(P_4) = (f_4)_4((P_0)_4 + N_4)$$

$$= (f_4)_4(P_0)_4 + (f_4^4)_4(P_0)_4 N_4 + o_4(N_4), \qquad (16)$$

$$(f_2)_4(Q_4)$$

$$= (f_2^4)_4(0)Q + o_4(Q_4) \quad (17)$$

Write $f_4(P_0) = f_{40}, f_4^{'}(P_0) = f_{40}^{'}$, and $f_2^{'}(0) = f_{20}^{'}$ and substitute (16) and (17) into (15) to obtain

$$(u_4)_4 + (u_a)_4 \approx -(\mu_4(a_4) + (f_{40})_4)u - (f_{40}^4)_4(\rho_0)_4(N_4) - (f_{20}^4)_4(\rho_0)_4(Q_4)$$

$$u_4(t,0) = \rho_4(t,0)(\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a) + (\rho_0)_4(a_4)da_4 - (\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a_4)da_4 + (h_4(Q_4) - 4)(\rho_0)_4(0) \quad (18)$$

Expanding

$$h_4(Q_4) = 1 + (h^4)_4(0)Q_4 + 0(Q_4) \equiv 1 + h_0^{'}Q_4 + 0Q_4$$

And substituting it into (48), we arrive at

$$u_4(t,0) \approx \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + (h_{a_4}^{'})(\rho_0)_4Q_4$$

The linearization is the following :

$$[u_4] + [u_4] = -(\mu_4(a_4) + [t_{40}]_4u_4 - [\alpha_4]_4(a_4)N_4 - [\alpha_2]_4(a_4)Q_4$$

$$u_4(t,0) = \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + [\xi_4]Q_4 \quad (19)$$

$$\frac{dQ_4}{dt} = -b_4Q_4$$

where

$$[\alpha_4]_4(a_4) = [t_4^4 a]_4[p, o]_4(a_4)d2(a_4) = [f_{20}^4]_4(\rho\, o\,)_4(a_4)\xi = [\mathrm{h\ o}^4]_4[\rho \mathrm{o}]_4(\mathrm{o}) \text{ and } \mathrm{b}_4 = \mathrm{b}_4 - g_4(P_4)$$

If $g_4(P_0) > b$, clearly, $(\rho_0)_4(a_4), 0)$ is unstable and the predator population is persistent. Assume $g_4(P_0)_4 < b$.

Then, $Q_4(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_4 + u_a = -(\mu_4(a_4) + (f_{40})_4)u_4 - (\alpha_4)_4(a_4)N_4$$

$$= -(\mu_4(a_4) + f_{40})u_4 - \alpha_4(a_4)\int_0^\infty u_4(t, a_4)da_4$$

$$u_4(t, 0) = \int_0^\infty \beta_4\,(a_4)u_4(t, a_4)da_4$$

Using the method developed by Cushing [2, 3], we suppose

$$u_4(t, a_4) = \tilde{\mathrm{u}}_4(a_4)^{e^{\gamma(t-a_4)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_4}\, da_4$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_4)$ gives,

$$\tilde{\mathrm{u}}_4(a_4) = e^{-\int_0^a \mu_4(\tau_4)d\tau - (f_{40})_4 a}\left(\tilde{\mathrm{u}}_4(0) - w\int_0^a e^{\int_0^\tau \mu_4(s_4)ds + (f_{40})}\tau_4\,(\alpha_4)_4(\tau_4)e^{\gamma_4\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_4(a_4)e^{-\gamma_4 a_4 - (f_{40})_4 a_4 - \int_0^a \mu_4(\tau_4)d\tau}\right)\tilde{\mathrm{u}}_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a(\alpha_4)_4(\tau_4)e^{-\gamma_4\tau_4}e^{(f_{40})_4\tau_4+\int_0^{\tau}\mu_4(s_4)ds}d\tau_4\,da_4$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}da\right)(\tilde{u}_4(0))_4$$

$$+\left(-4-\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a e^{(f_{40})_4\tau_4+\int_0^{\tau}\mu_4(s_4)ds+y\tau}\alpha_4(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(t_4)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_4\,(a_4)Y_4(\gamma_4,a_4)da_4\right)u_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)\int_0^a(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4=0,$$

(24)

$$\left(\int_0^{\infty}Y_4(\gamma_4,a_4)da_4\right)\tilde{u}_4(0)$$

$$+\left(-1-\int_0^{\infty}Y_4(\gamma_4,a_4)\int_0^a(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4\right)w=0.$$

Equqtion (20) has a solution of the form of $u_4(t,a_4) = u_4(a_4)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_4(0) \not\equiv 0$ and $\gamma, w$ in the complwx plane.

Let $I_4(\gamma_4,a_4) = \int_0^a (\rho_0)_4(\tau_4)Y^{-4}(\gamma_4,a_4)d\tau_4$. We obtain the characteristic equation

$$f_{40}^{l}{}_4\left(\int_0^\infty Y_4(\gamma_4,a_4)da_4\right)\left(\int_0^\infty \beta_4(a_4)Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$+\left(1-\int_0^\infty \beta_4(a_4)Y_4(\gamma_4,a_4)da\right)\left(1+(f_{40}^{l})_4\int_0^\infty Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$= 0. \qquad (25)$$

Since

$$(\rho_0)_4(a_4) = (\rho_0)_4(0)e^{-\int_0^a \mu_4(\tau_4)d\tau_4-(f_{40})_4a_4}$$

$$I_4(\gamma_4,a_4) = \int_0^a (\rho_0)_4(\tau_4)e^{\gamma_4\tau_4+(f_{40})_4\tau+\int_0^\tau \mu_4(s_4)ds_4}d\tau$$

$$= \frac{(\rho_0)_4(0)}{\gamma_4}(e^{\gamma_4a_4}-1)$$

Let

$$E_4(a_4) = e^{-(f_{40})_4a_4-\int_0^a \mu_4(\tau_4)d\tau_4}$$

Then $[Y_4]_4[\gamma_4]_4,a_4) = E_4(a_4)e^{-\gamma_4a_4}$. by a simple calculation,

$$\int_0^\infty \beta_4(a_4)Y_4(\gamma_4,a_4)I(\gamma_4,a_4)da = \frac{[\rho_0]_4(0)}{\gamma}\left(1-\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}\ da\right);$$

$$\int_0^\infty Y_4(\gamma_4,a_4)I(\gamma_4,a_4)da = \frac{[\rho_0]_4(0)}{\gamma_4}\left(\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}\ da_4\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_4(0)}{\gamma_4}[f_{40}^4]\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}da_4\ (\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4)$$

$$+\left(1-\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right).$$

$$\left(1+\frac{(\rho_0)_4(0)}{\gamma_4}(f_{40}^{'})_4\left(\int_0^\infty E_4(\,a_4)da_4-\int_0^\infty E_4(\,a_4)e^{-\gamma_4 a_4}da_4\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right)(1+\frac{(f_{40}^{'})_4}{\gamma_4}(\rho_0)_4(0)\int_0^\infty E_4(\,a_4)da_4)=0.$$

We need ũ(a) $\not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)da_4 = 1-\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4 \neq 0,$$

Since $\int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)\ \int_0^a \alpha_4\,(\tau_4)Y^{-4}(\gamma_4,\tau_4)d\tau_4 da_4 > 0$. Therefore, the characteristic solution is

$$\gamma_4 = -(f_{40}^{'})_4(\rho_0)_4 \int_0^{\infty} E_4(a_4)da_4.$$

Since $f_4'(P_0)_4 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_4), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_4), 0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_4 + (\mu_4(a_4) + [f_4]_4(P_4^*) + [f_2]_4(Q_4^*))P_4^* = 0, \quad (26)$$

$$p_4^*(0) = h_4(Q_4^*)\int_0^{\infty} \beta_4(a_4)p_4^*(a_J)p_4^*(a_4)da_4 \quad (27)$$

$$g_4(P_4^*) - b_4 = 0, \quad (28)$$

Where $P_4^* = \int_0^{\infty} p_4^*(a_4)\, da_4$

Equation (28) yields $P^* = g^{-4}(b)$,

From (26),

$$p_4^*(a) = p_4^*(0)e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}$$

Substituting into (27), we need to solve

$$\left(1 - h_4(Q_4^*)\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}da_4\right)p_4^*(0) = 0$$

For $Q_4^*$

Fix $p_4^*(0)$ and $P_4^*$ Let

$$H_4(z_4) = 1 - h_4(z_4)\int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}da_{,4}$$

And assume that

$$[k_3]_4\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(g_4^{-4}(b_4))a_4 - [k_2]_4 a_4}\, da_4 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_4(z_4)] > 0.$$

Clearly,

$$\frac{dH_4}{dz_4} = h_4(z)([f_2']_4(z_4))\int_0^\infty a_4\beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - f_4(P_4^*)a_4 - f_2(z_4)a_4}da_{\square}$$

$$-h'(z_4)\int_0^\infty \beta_4(a_4)\, e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}\quad da_4$$

the assumptions on $\left[f_2^4\right]_4 and\ [h^4]^4 dH_4/dz > 0$. Hence, if

$$H_4(0) = 1 - \int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4} da_4 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_4(0) < 0$ follows if $g_4^-(b_4) < [f_4^{-4}]_4(r_4)$, where $r_4$ is the solution of (12).

Indeed, since

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 e^{-y\frac{a}{4}} da_4}$$

Is a decreasing function of $s_4$ when $P_4^* = [g^{-4}]_4(b_4) < [f_4^{-4}]_4(r_4)$, that is $[f_4]_4(P_4^*) < r$, $R\left(f_4(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_4(g^{-4}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4\left(g_4^{-4}(b)\right)a_4} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_4^*(a_4), Q_4^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_4^*(a_4), Q_4^*)$ is more complicated. We again consider perturbations $m_4(t_4 a_4) = p_4(t_4, a_4) - p_4^*(a_4)$ and $W_4(t) = Q_4(t) - Q^*$.Then

$$M_4(t) \equiv \int_0^\infty m_4(t, a_4) da_4 = P_4(t) - P_4^*,$$

$$[m_4]_4 + [m_a]_4 = [p_4]_4 + [p_a]_4 - [p^*]_4(a_4) \approx -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4$$

$$-[f_{4*}^4]_4 p_4^*(a_4)M - [f_{2*}^4]_4 p_4^*(\mathrm{a})W_4,$$

$$m_4(t,0)p_4(t,0) - p_4^*(0) \approx h_* \int_0^\infty \beta_4(a_4)m_4(t,a_4)\,da + \frac{[h_*^4]_4}{[h_*]_4}\, p_4^*(0)W_4$$

$$\frac{dW_4}{dt} = \frac{dQ_4}{dt} \approx [g_*^4]_4 Q_4^* M_{,4}$$

Where $[f_{4*}]_4$ denoting $[f_4(P^*)]_4$, $[f_{4*}]_4$ denoting $[f_2(P^*)]_4$, $[f_{4*}^{'}]_4$ denoting $[f^{'}{}_4(P^*)]_4 [f_{2*}^{'}]_4$ denoting $[f_2]_4$ $(Q_4^*,$ $h_{*}^{'}$ denoting $h_*^{'}$ denoting $h_4^{'}$ $(Q_4^*)$, and $[g_*^{'}]_4$

Denoting $[g^{'}]_4(P_4^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_4]_4 + [m_a]_4 = -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4 - [f^{'}{}_{4*}]_4 [p_4^*](a_4)M_4 - [f_{2*}^{'}]_4 p_4^*(a_4)W_4$$

$$m_4(t,0) = [h_*]_4 \int_0^\infty \beta_4(a_4)m_4(t,a_4)da_4 + \frac{[h_*^4]_4}{[h^*]_4}[p^*]_4(0)\, W_{,4} \qquad (30)$$

$$\frac{dW_4}{dt} = [g_*^4]_4 [Q_4^*]M_4$$

This is equation is important in the following cases: problem of using observed correlations to infer causal relations is relevant to (e) a wide variety of scientific disciplines. Yet given correlations between just two classical variables, (eb) it is impossible to determine whether they arose from (e) a causal influence of one on the other or a common cause influencing both, unless one can implement (eb) a randomized intervention. **Katja Ried, Megan Agnew,**

**Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. Resch** here consider the problem of causal inference for (e) quantum variables. Authors introduce causal tomography, which unifies and generalizes (e&eb) conventional quantum tomography schemes to provide (eb) a complete solution to the causal inference problem using (e) a quantum analogue of a randomized trial. Authors furthermore show that, in contrast to the classical case, observed quantum correlations alone can sometimes provide (eb) a solution. They implement a quantum-optical experiment that allows (eb) to control (e&eb) the causal relation between two optical modes, and two measurement schemes -- one with (e&eb) and one without (e) randomization -- that extract this relation from (e) the observed correlations. Results show that entanglement and (e&eb) coherence, known to be central to (e) quantum information processing, also provide (eb) a quantum advantage for (e) causal inference. Subjects: Quantum Physics (quant-ph); Learning (cs.LG); General Relativity and Quantum Cosmology (gr-qc); Machine Learning (stat.ML) Journal reference: Nat Phys 11, 414-420 (2015) DOI: 10.1038/nphys3266 Cite as: arXiv: 1406.5036 [quant-ph] (or arXiv: 1406.5036v1 [quant-ph] for this version) **Inferring causal structure: a quantum advantage Katja Ried, Megan Agnew, Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. ReschSTOP D. T. Son, A. O. Starinets** review recent progress in applying the AdS/CFT correspondence to (e&eb) finite-temperature field theory. In particular, they show how the hydrodynamic behavior of field theory is reflected in (e&eb) the low-momentum limit of (e) correlation functions computed through (e&eb) a real-time AdS/CFT prescription, which they formulate. Authors also show how the hydrodynamic modes in field theory correspond to (e&eb) the low-lying quasinormal modes of (e) the AdS black p-brane metric. They provide a proof of the universality of (e) the viscosity/entropy ratio within (eb) a class of theories with (e&eb) gravity duals and formulate (eb) a viscosity bound conjecture. Possible implications for real systems are mentioned. Subjects: High Energy Physics - Theory (hep-th) Journal reference: Ann.Rev.Nucl.Part.Sci.57:95-118,2007 DOI: 10.1146/annurev.nucl.57.090506.123120Report number: INT PUB 07-02Cite as: arXiv:0704.0240 [hep-th]

(or arXiv:0704.0240v2 [hep-th] for this version) An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society **Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP** Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv:

1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version **Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP Viscosity, Black Holes, and Quantum Field Theory D. T. Son, A. O. Starinets STOP Authors** study information retrieval from evaporating black holes, assuming that the internal dynamics of a black hole is unitary and rapidly mixing, and assuming that the retriever has unlimited control over the emitted Hawking radiation. If the evaporation of the black hole has already proceeded past the ``half-way" point, where (e) half of the initial entropy has been radiated away, then (eb) additional quantum information deposited in (eb) the black hole is revealed in (eb) the Hawking radiation very rapidly. Information deposited prior to the half-way point remains concealed until (e) the half-way point, and then emerges (eb) quickly. These conclusions hold because (e) typical local quantum circuits are efficient encoders for (e) quantum error-correcting codes that nearly achieve (eb) the capacity of the quantum erasure channel. Estimate of a black hole's information retention time, based on (e) speculative dynamical assumptions, is just barely compatible with (e&eb) the black hole complementarity hypothesis. **Black holes as mirrors: quantum information in random subsystems Patrick Hayden1 and John Preskill2 Published 26 September 2007 Journal of High Energy Physics, Volume 2007, JHEP09 (2007) STOP** By analyzing the key properties of black holes from the point of (e) view of quantum information, authors derive a model-independent picture of (e) black hole quantum computing. It has been noticed that this picture exhibits (eb) striking similarities with (e&eb) quantum critical condensates, allowing (eb) the use of a common language to describe (eb) quantum computing in both systems. Authors analyze such quantum computing by allowing coupling to (e&eb) external modes, under the condition that the external influence must be (=) soft-enough in order not to (e) offset the basic properties of the system. Authors derive model-independent bounds on some crucial time-scales, such as (=) the times of gate operation, decoherence, maximal entanglement and total scrambling. They show that for black hole type quantum computers all these time-scales are of (e) the order of the

black hole half-life time. Furthermore, authors construct explicitly a set of Hamiltonians that generates (eb) a universal set of quantum gates for (e) the black hole type computer. They find that the gates work at maximal energy efficiency. Furthermore, authors establish a fundamental bound on the complexity of quantum circuits encoded on (eb&e) these systems, and characterize the unitary operations that are (=) implementable. It becomes apparent that the computational power is very limited due to (e) the fact that the black hole life-time is of the same order of (e) the gate operation time. As a consequence, it is impossible to retrieve its information, within (eb) the life-time of a black hole, by externally coupling to (e&eb) the black hole qubits. However, authors show that, in principle, coupling to (e&eb) some of the internal degrees of freedom allows (eb) acquiring knowledge about the micro-state. Still, due to the trivial complexity of operations that can be performed, (eb) there is no time advantage over the collection of Hawking radiation and subsequent decoding. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Report number: LMU-ASC 03/16, MPP-2016-1 Cite as: arXiv:1605.01407 [hep-th] (or arXiv:1605.01407v1 [hep-th] for this version) **Universality of Black Hole Quantum Computing Gia Dvali, Cesar Gomez, Dieter Lust, Yasser Omar, and Benedikt Richter (Submitted on 4 May 2016) STOP** In this addendum to [arXiv: 1402.5674] two points are discussed. In the first additional evidence is provided (eb) for a dual connection between the geometric length of an Einstein-Rosen bridge and (e&eb) the computational complexity of the quantum state of the dual CFT's. The relation between growth of complexity and (e&eb) Page's “Extreme Cosmic Censorship” principle is also remarked on. The second point involves a gedanken experiment in which Alice measures (eb) a complete set of commuting observables at her end of (e) an Einstein-Rosen bridge is discussed. An apparent paradox is resolved by appealing to (e) the properties of GHZ tripartite entanglement. **Addendum to computational complexity and black hole horizons Leonard Susskind* Version of Record online: 18 JAN 2016 DOI: 10.1002/prop.201500093 STOP** Analyzing some well established facts, authors give a model-independent

parameterization of black hole quantum computing in terms of (e&*eb) a set of macro and micro quantities and their relations. These include the relations between the extraordinarily-small energy gap of (e) black hole qubits and important time-scales of (e) information-processing, such as, scrambling time and Page's time. Authors then show, confirming and extending previous results, that other systems of nature with (e&eb) identical quantum informatics features are (=) attractive Bose-Einstein systems at (eb) the critical point of quantum phase transition. Here authors establish a complete isomorphy between (e&eb) the quantum computational properties of these two systems. In particular, they show that the quantum hair of a critical condensate is (=) strikingly similar to (=) the quantum hair of a black hole. Irrespectively whether one takes the similarity between the two systems as a remarkable coincidence or as (=) a sign of a deeper underlying connection, the following is evident. Black holes are not unique in their way of (e) quantum information processing and they can manufacture black hole based quantum computers in labs by (e) taking advantage of quantum criticality. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Atomic Physics (physics. Atom-ph); Quantum Physics (quant-ph) Cite as: arXiv: 1601.01329 [hep-th] (or arXiv: 1601.01329v1 [hep-th] for this version) **Black Hole Based Quantum Computing in Labs and in the Sky Gia Dvali, Mischa Panchenko (Submitted on 6 Jan 2016) STOP Gia Dvali, Alexander Gußmann** show that the existence of black holes with (e&eb) classical skyrmion hair invalidates (e) standard proofs that global charges, such as (=) the baryon number, cannot be conserved by (e) a black hole. By carefully analyzing the standard arguments based on (e) a Gedankenexperiment in which a black hole is seemingly-unable to (e) return the baryon number that it swallowed, authors identify inconsistencies in this reasoning, which does not (e) take into the account neither the existence of skyrmion black holes nor the baryon/skyrmion correspondence. Authors then perform a refined Gedankenexperiment by incorporating (e) the new knowledge and show that (eb) no contradiction with conservation of baryon number takes place at (eb) any stage of black hole evolution. Analysis also indicates (eb) no conflict

between semi-classical black holes and (e&eb) the existence of baryonic gauge interaction arbitrarily-weaker than gravity. Next, authors study classical cross sections of a minimally-coupled massless probe scalar field scattered by (e) _ a skyrmion black hole. Authors investigate how the skyrmion hair manifests itself by (e) comparing this cross section with (e&eb) the analogous cross section caused by (e) a Schwarzschild black hole which has (e) the same ADM mass as (=) the skyrmion black hole. Here authors find an order-one difference in the positions of (e) the characteristic peaks in (eb) the cross sections. The peaks are shifted to smaller scattering angles when (e) the skyrmion hair is present. This comes from the fact that the skyrmion hair changes (e&eb) the near horizon geometry of the black hole when compared to (e&eb) a Schwarzschild black hole with same ADM mass. Authors keep the study of this second aspect general so that (e) the qualitative results obtained can also be applied to(e&eb) black holes with classical hair of different kind. Subjects: High Energy Physics - Theory (hep-th); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU-ASC 18/16 Cite as: arXiv:1605.00543 [hep-th] (or arXiv:1605.00543v2 [hep-th] for this version) A preliminary discussion is given of (e) the prospects that gravitational-wave observations of (e) binary inspiral of black holes could reveal or constrain (e&eb) quantum modifications to (e) black hole dynamics, such as are required to (e) preserve postulates of quantum mechanics. Different proposals for such modifications are characterized by (e) different scales, and the size of these scales relative to (e&eb) those probed by observation of (e) inspiral signals is important in determining (eb) the feasibility of finding experimental signatures. Certain scenarios with strong quantum modifications in (eb) a region extending well outside the horizon are expected to modify (e&eb) classical evolution, and distort (e&eb) the near-peak gravitational wave signal, suggesting (eb) a search for departures from (e) waveforms predicted by general relativity. The near agreement of the GW150914 signals with (e&eb) such waveforms is discussed, and indicates (eb) constraints on some such scenarios. Important strategies for more precise future tests are 1) to develop more precise

predictions from (e) scenarios proposing quantum modifications, and 2) searching for observed deviations from (e) numerical relativity predictions via (e&eb) analysis of gravity wave data, particularly focussing on (e&eb) the signal region corresponding to plunge and merger. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1602.03622 [gr-qc] (or arXiv: 1602.03622v2 [gr-qc] for this version) **Gravitational wave tests of quantum modifications to black hole structure -- with post-GW150914 update Steven B. Giddings STOP** Direct LHC bounds on (e&eb) colored SUSY particles now corner naturalness more than the measured value of the Higgs mass does. Bounds on the gluino are of (e) particular importance, since it radiatively "sucks" up (e) the stop and Higgs-up soft masses. As a result, even models that easily accommodate a 125 GeV Higgs are almost as tuned as (=) the simplest version of SUSY, the MSSM: at best at the percent level. In this paper, authors further examine how current LHC results constrain (e) naturalness in three classes of models that may relax LHC bounds on (e&eb) sparticles: split families (e&eb), baryonic RPV (e&eb), and Dirac gauginos. In models of split families and bRPV, the bounds on the gluino are (=) only slightly reduced, resulting in (eb) a few percent tuning. In particular, having a natural spectrum in bRPV models typically implies (eb) that tops, Ws, and Zs are easily produced (eb) in the cascade decays of (e) squarks and gluinos. The resulting leptons and missing energy push (eb) the gluino mass limit above 1 TeV. Even when the gluino has a Dirac mass and does not contribute to the stop mass at one loop, tuning reappears in calculable models because there is no symmetry imposing the supersoft limit. Authors conclude that, even if sparticles are found at LHC-14, naturalness will not emerge triumphant. Subjects: High Energy Physics - Phenomenology (hep-ph)DOI:10.1007/JHEP03(2014)022 Cite as:rXiv:1309.3568 [hep-ph] (or arXiv:1309.3568v2 [hep-ph] for this version) **The Last Vestiges of Naturalness Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang, Ken Van Tilburg, Giovanni Villadoro STOP** The lack of evidence for new physics beyond (e) the standard model at the LHC points to (e) a paucity of new particles near the weak scale. This suggests that the weak scale is tuned and that supersymmetry, if present at

all, is realized (eb) at higher energies. The measured Higgs mass constrains (e) the scalar sparticles to be below 10^5 TeV, while (e) gauge coupling unification favors (eb) Higgsinos below 100 TeV. Nevertheless, in many models gaugino masses are (=) suppressed and remain within reach of the LHC. Tuning the weak scale and (e&eb) the renormalization group evolution of (e) the scalar masses constrain (e) Split model building. Due to the small gaugino masses, either the squarks or the up-higgs often run tachyonic; in the latter case, successful electroweak breaking requires heavy higgsinos near the scalar sparticles. We discuss the consequences of tuning the weak scale and the phenomenology of several models of Split supersymmetry including anomaly mediation, U(1)_(B-L) mediation, and Split gauge mediation. Subjects: High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Experiment (hep-ex) DOI: 10.1007/JHEP02 (2013)126 Cite as: arXiv: 1210.0555 [hep-ph] (or arXiv: 1210.0555v2 [hep-ph] for this version) **Mini-Split Asimina Arvanitaki, Nathaniel Craig, Savas Dimopoulos, Giovanni Villadoro STOP** Light states associated with (e&eb) the hierarchy problem affect (e&eb) the Higgs LHC production and decays. Authors illustrate this within the MSSM and two simple extensions applying (e&eb) the latest bounds from (e) LHC Higgs searches. Large deviations in the Higgs properties are expected in (eb) a natural SUSY spectrum. The discovery of a non-Standard-Model Higgs may signal (eb) the presence of light stops accessible at (eb) the LHC. Conversely, the more the Higgs is Standard-Model-like, (eb) the more tuned the theory becomes. Taking the ratio of different Higgs decay channels at the LHC cancels (er&eb) the leading (eb) QCD uncertainties and potentially improves (eb) the accuracy in Higgs coupling measurements to (e) the percent level. This may lead to (eb) the possibility of doing precision Higgs physics at (eb) the LHC. Finally, authors entertain the possibility that the ATLAS excess around (e&eb) 125 GeV persists with (e&eb) a Higgs production cross-section that is enhanced compared to (e&eb) the SM. This increase can only be accommodated in extensions of (e) the MSSM and it may suggest (eb) that stops lie below 400 GeV, likely within (eb) reach of next year's LHC run. Subjects: High Energy Physics - Phenomenology (hep-ph) DOI: 10.1007/JHEP02 (2012)144 Cite

as: arXiv: 1112.4835 [hep-ph] (or arXiv: 1112.4835v2 [hep-ph] for this version) **A Non Standard Model Higgs at the LHC as a Sign of Naturalness Asimina Arvanitaki, Giovanni Villadoro STOP** It has recently been suggested that the presence of a plenitude of light axions, an Axiverse, is (=) evidence for the extra dimensions of string theory. Authors discuss the observational consequences of (e) these axions on astrophysical black holes through (e&eb) the Penrose superradiance process. When an axion Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e&eb) the black hole "nucleus" forming (eb) a gravitational atom in the sky. The occupation number of superradiant atomic levels fed by (e) the energy and angular momentum of the black hole grows exponentially. The black hole spins down and an axion Bose-Einstein condensate cloud forms (eb) around it. When the attractive axion self-interactions become (=) stronger than the gravitational binding energy, (eb) the axion cloud collapses, a phenomenon known in condensed matter physics as "Bosenova". The existence of axions is first diagnosed by (e) gaps in the mass vs (e&eb) spin plot of astrophysical black holes. For young black holes the allowed values of (e) spin are quantized, giving rise to (eb) "Regge trajectories" inside the gap region. The axion cloud can also be observed directly either through (e&eb) precision mapping of the near horizon geometry or through (e&eb) gravitational waves coming from (e) the Bosenova explosion, as well as axion transitions and (e&eb) annihilations in (eb) the gravitational atom. Estimates suggest (eb) that these signals are detectable in (eb) upcoming experiments, such as Advanced LIGO, AGIS, and LISA. Current black hole spin measurements imply (eb) an upper bound on (e&eb) the QCD axion decay constant of (e) 2 x 10^17 GeV, while (e) Advanced LIGO can detect (e&eb) signals from (e) a QCD axion cloud with (e&eb) a decay constant as low as (=) the GUT scale. Authors finally discuss the possibility of observing the gamma-rays associated with the Bosenova explosion and, perhaps, the radio waves from axion-to-photon conversion for the QCD axion. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Journal

reference:Phys.Rev.D83:044026, 2011 DOI: 10.1103/PhysRevD.83.044026 Cite as: arXiv: 1004.3558 [hep-th] (or arXiv: 1004.3558v2 [hep-th] for this version) **Exploring the String Axiverse with Precision Black Hole Physics Asimina Arvanitaki, Sergei Dubovsky (Submitted on 20 Apr 2010 (v1), last revised 16 Aug 2011 (this version, v2)) STOP Skyrmion Black Hole Hair: Conservation of Baryon Number by Black Holes and Observable Manifestations Gia Dvali, Alexander Gußmann STOP Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust** discuss BMS supertranslations both at (eb) null-infinity and on the horizon for the case of the Schwarzschild black hole. Authors show that both kinds of supertranslations lead to (eb) infinitely many gapless physical excitations. On this basis authors construct a quotient algebra using (e) suited superpositions of (e) both kinds of transformations which cannot be compensated by (e) an ordinary BMS-supertranslation and therefore are (=) intrinsically due to (e) the presence of an event horizon. Authors show that these quotient transformations are (=) physical and generate (eb) gapless excitations on the horizon that can account for (e) the gravitational hair as well as for (e) the black hole entropy. They identify the physics of these modes as associated with (e&eb) Bogolioubov-Goldstone modes due to quantum criticality. Classically the number of these gapless modes is (=) infinite. However, authors show that due to (e) quantum criticality the actual amount of information-carriers becomes (=) finite and consistent with Bekenstein entropy. Although they only consider the case of Schwarzschild geometry, the arguments are extendable to(e&eb) arbitrary space-times containing event horizons. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU--ASC 05/16, MPP--2016—3 Cite as: arXiv: 1601.03725 [hep-th] (or arXiv: 1601.03725v1 [hep-th] for this version) **Gravitational Black Hole Hair from Event Horizon Supertranslations Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust STOP Gia Dvali, Cesar Gomez, Nico Wintergerst** suggest that holography can be formulated in terms of (e&eb) the information capacity of (e) the St\"uckelberg degrees of freedom that maintain

gauge invariance of (e) the theory in the presence of (e) an information boundary. These St\"uckelbergs act as (e&eb) qubits that account for a certain fraction of (e) quantum information. Their information capacity is measured by (e) the ratio of the inverse St\"uckelberg energy gap to (e) the size of the system. Systems with (e&eb) the smallest gap are (=) maximally holographic. For massless gauge systems this information measure is (=) universally equal to (=) the inverse coupling evaluated at the systems' length scale. In this language it becomes very transparent why the St\"uckelberg information capacity of black holes saturates (e full) the Bekenstein bound and accounts for (e) the entire information of the system. The physical reason is that the strength of quantum interaction is bounded from below by the gravitational coupling, which scales as (=) area. Observing the striking similarity between the scalings of the energy gap of the boundary St\"uckelberg modes and (e&eb): (AC=DC=0) the Bogoliubov modes of critical many-body systems, authors establish a connection between holography and (e&eb) quantum criticality through (e&eb) the correspondence between these modes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1511.03525 [hep-th] (or arXiv: 1511.03525v1 [hep-th] for this version) **Stückelberg Formulation of Holography Gia Dvali, Cesar Gomez, Nico Wintergerst STOP B**lack hole entropy, denoted by N, in (semi)classical limit is (=) infinite. This scaling reveals (eb) a very important information about the qubit degrees of (e) freedom that carry black hole entropy. Namely, the multiplicity of qubits scales as N, whereas (e) their energy gap and their coupling as 1/N. Such a behavior is indeed exhibited by (e) Bogoliubov-Goldstone degrees of freedom of (e) a quantum-critical state of N soft gravitons (a condensate or a coherent state) describing (eb) the black hole quantum portrait. They can be viewed as (=) the Goldstone modes of (e) a broken symmetry acting on (e&eb) the graviton condensate. In this picture Minkowski space naturally emerges as (=) a coherent state of infinite-N gravitons of (e) infinite wavelength and it carries (e&eb) an infinite entropy. In this paper authors **Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015)** ask what is the geometric meaning (if any) of (e) the classical limit of this symmetry. Authors argue that the infinite-

N limit of Bogoliubov-Goldstone modes of critical graviton condensate is described by (e) recently-discussed classical BMS super-translations broken by the black hole geometry. However, the full black hole information can only be recovered for (e) finite N, since the recovery time becomes (=) infinite in classical limit in which N is (=) infinite. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) DOI: 10.1016/j.physletb.2015.11.073 Report number: MPP-2015-215, LMU-ASC 59/15 Cite as: arXiv: 1509.02114 [hep-th] (or arXiv: 1509.02114v1 [hep-th] for this version) **Classical Limit of Black Hole Quantum N-Portrait and BMS Symmetry Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015) STOP** Recent ideas about understanding physics of black hole information-processing in terms of (e&eb) quantum criticality allow us to (eb) implement black hole mechanisms of quantum computing within (eb) critical Bose-Einstein systems. The generic feature, uncovered both by analytic and numeric studies, is (=) the emergence at the critical point of (e) gapless weakly-interacting modes, which act as (=) qubits for information-storage at a very low energy cost. These modes can be effectively described in terms of (e&eb) either Bogoliubov or Goldstone degrees of freedom. The ground-state at the critical point is (=) maximally entangled and far from being classical. We confirm this near-critical behavior by a new analytic method. **Gia Dvali, Mischa Panchenko** compute growth of entanglement and show its consistency with (e&eb) black hole type behavior. On the other hand, in the over-critical regime the system develops (eb) a Lyapunov exponent and scrambles quantum information very fast. By, manipulating (e&eb) the system parameters externally, authors can put it in and out of various regimes and in this way control (e&eb) the sequence of information storage and processing. By using gapless Bogoliubov modes as (=) control qubits they design some simple logic gates. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Cite as: arXiv: 1507.08952 [hep-th] (or arXiv: 1507.08952v1 [hep-th] for this version) **Black Hole Type Quantum**

**Computing in Critical Bose-Einstein Systems Gia Dvali, Mischa Panchenko (Submitted on 31 Jul 2015) STOP Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst** establish a fundamental connection between quantum criticality of a many-body system, such as Bose-Einstein condensates, and (e&eb) its capacity of information-storage and processing. For deriving the effective theory of modes in the vicinity of the quantum critical point they develop a new method by mapping a Bose-Einstein condensate of N-particles onto (e&eb) a sigma model with a continuous global (pseudo)symmetry that mixes (e&eb) bosons of different momenta. The Bogolyubov modes of the condensate are mapped onto (e&eb) the Goldstone modes of the sigma model, which become (=) gapless at the critical point. These gapless Goldstone modes are (=) the quantum carriers of information and entropy. Analyzing their effective theory, authors observe the information-processing properties strikingly similar to (=) the ones predicted by the black hole portrait. The energy cost per qubit of information-storage vanishes in (eb) the large-N limit and the total information-storage capacity increases with (e&eb) N either exponentially or as a power law. The longevity of information-storage also increases with (e&eb) N, whereas (e) the scrambling time in (eb) the over-critical regime is controlled by (e&eb) the Lyapunov exponent and scales logarithmically with (e&eb) N. This connection reveals (eb) that the origin of black hole information storage lies in (eb) the quantum criticality of the graviton Bose-gas, and that much simpler systems that can be manufactured in table-top experiments can exhibit (eb) very similar information-processing dynamics. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 92, 125002 (2015) DOI: 10.1103/PhysRevD.92.125002 Cite as: arXiv:1507.02948 [hep-th] (or arXiv:1507.02948v1 [hep-th] for this version) **Nambu-Goldstone Effective Theory of Information at Quantum Criticality Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst STOP Gia Dvali, Cesar Gomez** show that a positive cosmological constant is (=) incompatible with the quantum-corpuscular resolution of de Sitter metric in form of (e) a coherent state. The reason is very general and is due to (e) the

quantum self-destruction of (e) the coherent state because of (e) the scattering of constituent graviton quanta. This process creates an irreversible quantum clock, which precludes (e) eternal de Sitter. It also eliminates (e) the possibility of Boltzmann brains and Poincare recurrences. This effect is expected to be (=) part of any microscopic theory that takes into account (e) the quantum corpuscular structure of the cosmological background. This observation puts the cosmological constant problem in a very different light, promoting it, from a naturalness problem, into (e&eb) a question of quantum consistency. We are learning that quantum gravity cannot tolerate (e) exceedingly-classical sources. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1412.8077 [hep-th]

(or arXiv:1412.8077v1 [hep-th] for this version) **Quantum Exclusion of Positive Cosmological Constant? Gia Dvali, Cesar Gomez (Submitted on 27 Dec 2014) STOP G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger** establish a connection between the ultra-Planckian scattering amplitudes in field and string theory and (e&eb) unitarization by black hole formation in these scattering processes. **Using as a guideline an explicit microscopic theory in which the black hole represents (eb) a bound-state of many soft gravitons at (eb) the quantum critical point, they were able to identify and compute (e&eb) a set of perturbative amplitudes relevant for (e) black hole formation. These are the tree-level N-graviton scattering S-matrix elements in (eb) a kinematical regime (called classicalization limit) where (e) the two incoming ultra-Planckian gravitons produce (eb) a large number N of soft gravitons. Authors compute these amplitudes by using (e) the Kawai-Lewellen-Tye relations, as well as scattering equations and (e&eb) string theory techniques. Authors discover that this limit reveals (eb) the key features of the microscopic corpuscular black hole N-portrait. In particular, the perturbative suppression factor of a N-graviton final state, derived from (e) the amplitude, matches the non-perturbative black hole entropy when (e) N reaches the quantum criticality value, whereas final states with different value of N are (=) either suppressed or excluded by non-perturbative**

**corpuscular physics. Thus they identify the microscopic reason behind the black hole dominance over other final states including (e) non-black hole classical object. In the parameterization of the classicalization limit the scattering equations can be solved exactly allowing (eb) us to obtain closed expressions for the high-energy limit of the open and closed superstring tree-level scattering amplitudes for a generic number N of external legs. Authors demonstrate matching and complementarity between the string theory and (e&eb) field theory** in different large-s and large-N regimes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1016/j.nuclphysb.2015.02.004 Report number: MPP-2014-320, LMU-ASC 52/14 Cite as: arXiv:1409.7405 [hep-th] (or arXiv:1409.7405v2 [hep-th] for this version) **Black Hole Formation and Classicalization in Ultra-Planckian 2 -> N Scattering G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger STOP** Gravitational backgrounds, such as black holes, AdS, de Sitter and inflationary universes, should be viewed as (=) composite of N soft constituent gravitons. It then follows that such systems are close to (e) quantum criticality of graviton Bose-gas to Bose-liquid transition. Generic properties of the ordinary metric description, including (e) geodesic motion or particle-creation in the background metric, emerge as (=) the large-N limit of quantum scattering of (e) constituent longitudinal gravitons. **Gia Dvali, Cesar Gomez** show that this picture correctly accounts for (e) physics of large and small black holes in AdS, as well as reproduces (eb) well-known inflationary predictions for cosmological parameters. However, it anticipates new effects not captured by (e) the standard semi-classical treatment. In particular, **Gia Dvali, Cesar Gomez** predict observable corrections that are sensitive to (e&eb) the inflationary history way beyond last 60 e-foldings. They derive an absolute upper bound on the number of e-foldings, beyond (e) which neither de Sitter nor inflationary Universe can be approximated by (e) a semi-classical metric. However, they could in principle persist in a new type of (e) it quantum eternity state. Authors discuss implications of this phenomenon for (e&eb) the cosmological constant problem. Subjects: High Energy Physics - Theory (hep-th); Cosmology and

Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Quantum Physics (quant-ph) DOI: 10.1088/1475-7516/2014/01/023 Cite as: arXiv: 1312.4795 [hep-th] (or arXiv:1312.4795v1 [hep-th] for this version) **Quantum Compositeness of Gravity: Black Holes, AdS and Inflation Gia Dvali, Cesar Gomez STOP** As quantum information science approaches the goal of constructing quantum computers, understanding loss of information through (e&eb) decoherence becomes (=) increasingly important. The information about a system that can be obtained from (e) its environment can facilitate quantum control and error correction. Moreover, observers gain (eb) most of their information indirectly, by (e) monitoring (primarily photon) environments of (e) the "objects of interest." Exactly how this information is inscribed in the environment is essential for (e) the emergence of "the classical" from (e) the quantum substrate. In this paper, **Robin Blume-Kohout, Wojciech H. Zurek** examine how many-qubit (or many-spin) environments can store (e) information about a single system. The information lost to (e) the environment can be stored redundantly, or it can be encoded in (eb) entangled modes of the environment. Authors go on to show that randomly chosen states of the environment almost always encode (e) the information so that (e) an observer must capture (e) a majority of the environment to deduce (eb) the system's state. Conversely, in the states produced by (e) a typical decoherence process, information about a particular observable of the system is stored redundantly. This selective proliferation of "the fittest information" (known as Quantum Darwinism) plays a key role in (eb) choosing the preferred, effectively classical observables of (e) macroscopic systems. The developing appreciation that the environment functions not just as (=) a garbage dump, but as a communication channel, is extending understanding of the environment's role in (e&eb) the quantum-classical transition beyond (e) the traditional paradigm of decoherence. Keywords Decoherence quantum Darwinism quantum information quantum-classical transition **Foundations of Physics November 2005, Volume 35, Issue 11, pp 1857-1876 A Simple Example of "Quantum Darwinism": Redundant Information Storage in Many-Spin Environments Robin Blume-Kohout, Wojciech H. Zurek STOP Harold Ollivier, David**

**Poulin, and Wojciech H. Zurek** study the emergence of objective properties in open quantum systems. In our analysis, the environment is promoted from a passive role of a reservoir selectively destroying quantum coherence to an active role of amplifier selectively **proliferating information about the system. We show that only preferred pointer states of the system can leave a redundant and therefore easily detectable imprint on the environment. Observers who—as is almost always the case—discover the state of the system indirectly (by probing a fraction of its environment) will find out only about the corresponding pointer observable. Many observers can act in this fashion independently and without perturbing the system. They will agree about its state. In this operational sense, preferred pointer states exist objectively. Received 24 July 2003 DOI:http://dx.doi.org/10.1103/PhysRevLett.93.220401 Objective Properties from Subjective Quantum States: Environment as a Witness Harold Ollivier, David Poulin, and Wojciech H. Zurek Phys. Rev. Lett. 93, 220401 – Published 22 November 2004 STOP** Quantum Darwinism recognizes that we—the observers—acquire our information about the "systems of interest" indirectly from their imprints on the environment. **Michael Zwolak, H. T. Quan, and Wojciech H. Zurek** show that information about a system can be acquired from a mixed-state, or hazy, environment, but the storage capacity of an environment fragment is suppressed by its initial entropy. In the case of good decoherence, the mutual information between the system and the fragment is given solely by the fragment's entropy increase. For fairly mixed environments, this means a reduction by a factor 1−h, where h is the haziness of the environment, i.e., the initial entropy of an environment qubit. Thus, even such hazy environments eventually reveal the state of the system, although now the intercepted environment fragment must be larger by ~ (1−h)−1 to gain the same information about the system. DOI:http://dx.doi.org/10.1103/PhysRevLett.103.110402 **Quantum Darwinism in a Mixed Environment Michael Zwolak, H. T. Quan, and Wojciech H. Zurek Phys. Rev. Lett. 103, 110402 – Published 8 September 2009** We study quantum Darwinism—the redundant recording of information about the preferred states of a decohering system by its environment—for an object illuminated by a

blackbody. In the cases of point-source and isotropic illumination, we calculate the quantum mutual information between the object and its photon environment. We demonstrate that this realistic model exhibits fast and extensive proliferation of information about the object into the environment and results in redundancies orders of magnitude larger than the exactly soluble models considered to date. DOI:http://dx.doi.org/10.1103/PhysRevLett.105.020404 **Quantum Darwinism in an Everyday Environment: Huge Redundancy in Scattered Photons C. Jess Riedel and Wojciech H. Zurek Phys. Rev. Lett. 105, 020404 – Published 9 July 2010 STOP** Measurements transfer information about a system to the apparatus and then, further on, to observers and (often inadvertently) to the environment. I show that even imperfect copying essential in such situations restricts possible unperturbed outcomes to an orthogonal subset of all possible states of the system, thus breaking the unitary symmetry of its Hilbert space implied by the quantum superposition principle. Preferred outcome states emerge as a result. They provide a framework for "wave-packet collapse," designating terminal points of quantum jumps and defining the measured observable by specifying its eigenstates. In quantum Darwinism, they are the progenitors of multiple copies spread throughout the environment—the fittest quantum states that not only survive decoherence, but subvert the environment into carrying information about them—into becoming a witness. DOI:http://dx.doi.org/10.1103/PhysRevA.76.052110 **Quantum origin of quantum jumps: Breaking of unitary symmetry induced by information transfer in the transition from quantum to classical Wojciech Hubert Zurek Phys. Rev. A 76, 052110 – Published 16 November 2007 STOP** The Born rule, a foundational axiom used to deduce probabilities of events from **wavefunctions, is indispensable in the everyday practice of quantum physics. It is also key in the quest to reconcile the ostensibly inconsistent laws of the quantum and classical realms, as it confers physical significance to reduced density matrices, the essential tools of decoherence theory. Following Bohr's Copenhagen interpretation, textbooks postulate the Born rule outright. However, recent attempts to derive it from (e) other quantum principles have been successful, holding promise for (e)**

**simplifying and clarifying the quantum foundational bedrock. A major family of derivations is based on envariance, a recently discovered symmetry of (e) entangled quantum states. Here, authors identify and experimentally test three premises central to (e) these envariance-based derivations, thus demonstrating, (eb) in the microworld, the symmetries from (e) which the Born rule is derived. Further, they demonstrate envariance in a purely local quantum system, showing (eb) its independence from relativistic causality.Subjects: Quantum Physics (quant-ph)Cite as:arXiv:1604.01471 [quant-ph] (or arXiv:1604.01471v1 [quant-ph] for this version) Quantum probabilities from quantum entanglement: Experimentally unpacking the Born rule Jérémie Harris, Frédéric Bouchard, Enrico Santamato, Wojciech H. Zurek, Robert W. Boyd, Ebrahim Karimi**

Integrating $[m_4]_4(t, a_4)$ with respect to a from 0 to ∞, we have

$$\frac{dM_4(t)}{dt} = \int_0^\infty [mt]_4(t, a_4)da_4$$

$$= -\int_0^\infty [m_a]_4\,(t, a_4)da_4 - \int_0^\infty \mu(a_4)m(t, a_4)da_4$$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f^{'}{}_{4*}]_4 \int_0^\infty p_4^*(a_4)da_4\right)M_4 - [f_{2*}^4]_4 \int_0^\infty p_4^*(a_4)da_4 W_4$$

$= m_4(t, 0) - \int_0^\infty \mu_4(a_4)\, m_4(t, a_4)da_4$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}^{i}]_4 P_4^*\right)M_4 - [f_{2*}^{i}]_4 P_4^* W_{\cdot 4}$$

From (30), we arrive at the equations

$$\frac{dW_4(t)}{dt} = \int_0^{\infty} ([h_*]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4$$

$$-\left([f_{4*}] + [f_{2*}]_4 + \left[f_{4*}^{'}\right]_4 P_4^*\right) M_4 + \left(\frac{\left[h_*^{'}\right]_4}{[h_*]^4}\, P_4^*(0) - f_{2*}^{'} P_4^*\right) W_4,$$

(34)

$$\frac{dW_4(t)}{dt}\left[g_{*}^{'}\right]_4 Q_4^*\, M_{.4}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{r*}]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4 = ([h_*]_4\beta_4(\theta_4) - \mu_4(\theta_4))M_{,4}$$

Where $0< \theta_4 = \theta_4(t) < \infty$.

And if $a_4^2\,[a_4^2]_4<4\,[a_2]_4<1\, a_2 e$,

$\mathrm{Re}\,\lambda_4\left[\lambda_4\right]_4 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_4 = -\frac{4}{2}a_4[a2\ ]_4 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_*\,[h_*]_4\beta_4(a_4) - \mu_4(a_4)\right|$. Then$|n_4(t)|$ ⊠⊿∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_4^2 <4\, a_2 e$ and $[a_4] > 2\Delta$,or $[a_4^2]$ ⊠⊿$[a_2]e$ and $(a_4 - \sqrt{(a_4^2)_4 - 4(a_2)_4 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+4}\right)$ and $\beta_4(a_4) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4) + d\tau_4 da} = \frac{4}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_4 + [\mu_0]$ 0.5. Assume$[f_4](x_4)= K_4 x_4$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_4(x_4)= [0x_4]/(1 + x_4)$; the functional responses $[[f_2](x) = [0]\, x_4/(1 + x_4)$and $h_4(x_4)=(1+\omega_4 x_4)/(1 + x_4)$(See Holling [7] and Rashevsky [13])

Equation (44),now has the form of

$$\frac{m_4 + [\mu_0] + [K][P_0])}{[m_4]_4 + (\mu_0)_4 + [K][P_0])\frac{3}{4}\,[K]_4[P_0]\frac{3}{4}} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(4 + P_0) < b_4$ i e $[P_0] < b/(\delta_4 - b_4)$then the equilibrium$(P_0)$ $[P_0]_4(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_4$

$[f_{4*}] = K_4\, P^*\, [f_{2*}] = \frac{\sigma Q_4^*}{4+Q_4^*} \quad [h_*] = \frac{4+\omega_4 Q_4^*}{4+Q_4^*}$

$[f_{4*}] = K_4\, [f_{2*}^4] = \frac{\sigma_4}{(4+Q_4^*)^2}\, [g_*^4] = \frac{\sigma_4}{(4+P_4^*)^2} \quad [h_*^4] = \frac{\omega_4 - 4}{(4+Q_4^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0$, K, b, $\delta$, and $\omega$ need to satisfy

$$\frac{\omega_4}{(m + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma)^2} + \frac{2\omega_4}{([m_4]_4 + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma_4)^3} < 1$$

And

$$\frac{1}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^2}+\frac{2\omega_4}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4}+(m_4+[\mu_0]_4+\frac{K_4b_4}{\delta_4-b_4}+\sigma_4)\,Q_4^*)^3$$

$$=m_4+[\mu_0]_4+3+\frac{K_4b_4}{\delta_4-b_4}(4+Q^*)(1+\omega_4Q_4^*$$

Now

$$\left[a_4^2\right]-1\left[a_2\right]e=\left(\frac{2K_4b_4}{\delta_4-b_4}\right)^2+\frac{1K_4b_4\sigma_4\,Q_4^*}{(\delta_4-b_4)(1+Q_4^*)}+\frac{\sigma_4^2Q_4^{*2}}{(1+Q_4^*)^2}$$

$$\frac{1\,b_4\sigma_4\,(\delta_4-b)Q^*}{\delta_4(1+Q_4^*)^2}-\frac{1\,(\delta_4-b_4)^2(1-\omega_4)Q_4^*}{\delta_4(1+\omega_4Q_4^*)(1+Q_4^*)}P^*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_4\,ae_4^{mu}\;[\mu_0]_4-\frac{a_4}{a_4+4}$$

Hence if $[a_4^2>4(a_2)_4e$ and

$$[a_4] = \frac{2K_4 b_4}{\delta_4 - b_4} + \frac{\sigma_4 \, Q_4^*}{1 + Q_4^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_4]_4 = [f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}] P_4^*$$

$$[a_2]_4 = \left[f_{2*}^4\right] P_4^* - \frac{h_*^4}{h_*} P^*(0)$$

$$e = [g_*^4]_4 \, Q_4^*$$

$$\eta_4(t) = \left[h_*\right]_4 \beta_4 \left(\theta_4(t)\right) - \mu_4(\theta(t)),$$

The system(34)can be written as the following system:

$$[M(t)]\left[a_4\right] M_4 - [a_2]_4 W_4 - [n(t)] M_4$$

$$w_4(t) = eM_4$$

Write

$$A = \begin{pmatrix} -[a_4]_4 - [a_2]_4 & \\ e & 0 \end{pmatrix}$$

$$B = (t) = \begin{pmatrix} [\eta(t)]_4 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_4 \\ W_4 \end{pmatrix}$$

The solution of (32)is

$$X_4(t) = e^{A,t}\,[X]_4(0) + \int_0 e^{At}$$

**Some of the other conceptual and thematic and discursive factorial representationalities to which the model is applicable are:every one-dimensional quantum mechanical Hamiltonian H1 can have (e) a partner H2 such that H1 and H2 taken together may be viewed as (=) the components of a supersymmetric Hamiltonian. The term 'supersymmetric Hamiltonian' is taken to mean (eb) a Hamiltonian defined in terms of (e&eb)charges that obey (eb)the same algebra as that of (e) the generators of supersymmetry in field theory. The consequences of this symmetry for (er) the spectra of H1 and H2 are explored. It is shown how the supersymmetric pairing may be utilised to eliminate (e) the ground state of H1, or add a state below the ground state of H1 or maintain the spectrum of H1. It is also explicitly demonstrated that the supersymmetric pairing may be used to generate (eb) a class of anharmonic potentials with (e&eb) exactly specified spectra. The complete spectrum of an anharmonic potential so generated consists of (e) all the eigenstates of the simple harmonic oscillator and, in addition, a ground state at a specified energy E which lies arbitrarily below (eb) the E=1/2 ground state of the harmonic oscillator. Supersymmetric quantum mechanics of one-dimensional systems C V Sukumar Journal of Physics A: Mathematical and General, Volume 18, Number 15 STOP Authors examine the properties and consequences of (e&eb) pseudo-supersymmetry for (e) quantum systems admitting (e) a pseudo-Hermitian Hamiltonian. Authors explore the Witten index of (e) pseudo-**

**supersymmetry and show that every pair of diagonalizable (not necessarily Hermitian) Hamiltonians with (e&eb) discrete spectra and (e&eb) real or complex-conjugate pairs of (e) eigenvalues are (=) isospectral and have (e) identical degeneracy structure except (e) perhaps for the zero eigenvalue if and only if (e) they are pseudo-supersymmetric partners. This implies (eb) that pseudo-supersymmetry is (=) the basic framework for (e) generating non-Hermitian PT-symmetric and (e&eb) non-PT-symmetric Hamiltonians with (e&eb) a real spectrum via (e&eb) a Darboux transformation, and shows (eb) that every diagonalizable Hamiltonian H with(e&eb) a discrete spectrum and (e&eb) real or complex-conjugate pairs of (e) eigenvalues may be factored as (=) H=L♯L where L is (=) a linear operator with pseudo-adjoint L♯. In particular, this factorization applies to (e&eb) PT-symmetric and Hermitian Hamiltonians. The non-degenerate two-level systems provide (eb) a class of Hamiltonians that are pseudo-Hermitian. Authors demonstrate the implications of general results for (e&eb) this class in some detail. Copyright © 2002 Elsevier Science B.V. All rights reserved. Nuclear Physics B Volume 640, Issue 3, 30 September 2002, Pages 419–434Pseudo-supersymmetric quantum mechanics and isospectral pseudo-Hermitian Hamiltonians Ali Mostafazadeh STOP Michele Caselle, Alessandro Nada, Marco Panero present a high-precision lattice calculation of (e) the equation of state in the confining phase of (e) SU(2) Yang-Mills theory. Authors show that the results are described very well by (e) a gas of massive, non-interacting glueballs, provided (e) one assumes an exponentially growing Hagedorn spectrum. The latter can be derived within (eb) an effective bosonic closed-string model, leading to (eb) a parameter-free theoretical prediction, which is in perfect agreement with (=) proposed lattice results. Furthermore, when applied to (e&eb) SU (3) Yang-Mills theory, this effective model accurately describes (eb) the lattice results reported by Bors\'anyi et al. in JHEP 07 (2012) 056. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 07 (2015) 143 DOI: 10.1007/JHEP07 (2015)143 Cite as: arXiv:**

**1505.01106 [hep-lat] (or arXiv: 1505.01106v2 [hep-lat] for this version) Hagedorn spectrum and thermodynamics of SU (2) and SU (3) Yang-Mills theories Michele Caselle, Alessandro Nada, Marco Panero STOP B. de Wit J. Hoppe, H. Nicolai study the quantum-mechanical properties of (e) a supermembrane and examine the nature of (e) its ground state. A supersymmetric gauge theory of (e) area-preserving transformations provides (eb) a convenient framework for this study. The supermembrane can be viewed as (=) a limiting case of a class of models in supersymmetric quantum mechanics. Its mass does not depend on (=) the zero modes and vanishes only if (e) the wave function is (=) a singlet under (e&eb) supersymmetry transformations of (e) the nonzero modes. Authors exhibit (eb) the complexity of (e) the supermembrane ground state and (e&eb) examine various truncations of these models. None of these truncations has massless states. Copyright © 1988 Published by Elsevier B.V. Nuclear Physics B Volume 305, Issue 4, 28 December 1988: Pages 545-581: On the quantum mechanics of supermembranes Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. B. de Wit J. Hoppe, H. Nicolai STOP In this note we explain how world-volume geometries of (e) D-branes can be reconstructed within (eb) the microscopic framework where (e) D-branes are described through (e&eb) boundary conformal field theory. Authors extract the (non-commutative) world-volume algebras from (e) the operator product expansions of (e) open string vertex operators. For branes in a flat background with (e&eb) constant non-vanishing B-field, the operator products are computed perturbatively to (e) all orders in the field strength. Resulting series coincides with (=) Kontsevich's presentation of the Moyal product. After extending these considerations to (e) fermionic fields authors conclude with some remarks on the generalization of approach to (e&eb) curved backgrounds. D-branes and deformation quantization Volker Schomerus1,2 Published 25 July 1999 Journal of High Energy Physics, Volume 1999, JHEP06(1999)STOP Authors employ the path integral approach developed in Gangopadhyay S. and Scholtz F. E., Phys. Rev.**

**Lett., 102 (2009) 241602 to discuss the (generalized) harmonic oscillator in (eb) a noncommutative plane. The action for (e) this system is derived in (e&eb) the coherent state basis with (e&eb) additional degrees of freedom. From this the action in the coherent state basis without (e) any additional degrees of freedom is obtained. This gives (eb) the ground-state spectrum of the system. Authors then employ the exact renormalization group approach to show (eb) that equivalence can be constructed between this (noncommutative) system and (e&eb) a commutative system Path integral action and exact renormalization group dualities for quantum systems in noncommutative plane Sunandan Gangopadhyay Published 25 June 2015 • Copyright © EPLA, 2015 STOP Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini present a lattice study of (e)the equation of state in Yang-Mills theory based on (e) the exceptional G(2) gauge group. As is well-known, at zero temperature this theory shares (e&eb) many qualitative features with real-world QCD, including (e) the absence of colored states in the spectrum and dynamical string breaking at (e) large distances. In agreement with previous works, authors show that at finite temperature this theory features (eb) a first-order deconfining phase transition, whose nature can be studied by (e) a semi-classical computation. Authors also show that the equilibrium thermodynamic observables in the deconfined phase bear striking quantitative similarities with (=) those found in SU (N) gauge theories: in particular, these quantities exhibit (eb) nearly perfect proportionality to the number of gluon degrees of freedom, and (e&eb) the trace anomaly reveals (eb) a characteristic quadratic dependence on (e&eb) the temperature, also observed in SU (N) Yang-Mills theories (both in four and in three spacetime dimensions). Authors compare our lattice data with (e&eb) analytical predictions from (e) effective models, and discuss their implications for (e) the deconfinement mechanism and (e&eb) high-temperature properties of (e) strongly interacting, non-supersymmetric gauge theories. Results give (eb) strong evidence for (e) the conjecture that the thermal deconfining transition is governed by (e&eb) a universal mechanism, common to (e) all simple gauge groups.**

**Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 03 (2015) 057 DOI: 10.1007/JHEP03 (2015)057 Report number: DESY-14-146; IFT-UAM/CSIC-14-076 Cite as: arXiv: 1409.8305 [hep-lat] (or arXiv: 1409.8305v2 [hep-lat] for this version Exceptional thermodynamics: The equation of state of G(2) gauge theory Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini STOP In U (1) lattice gauge theory in three spacetime dimensions, the problem of confinement can be studied analytically in (e&eb) a semi-classical approach, in terms of (e&eb) a gas of monopoles with (e&eb) Coulomb-like interactions. In addition, this theory can be mapped to (e&eb) a spin model via (e&eb) an exact duality transformation, which allows (eb) one to perform high-precision numerical studies of (e) the confining potential. Taking advantage of these properties, authors carried out an accurate investigation of (e) the effective string describing (eb) the low-energy properties of flux tubes in (eb) this confining gauge theory. Authors found striking deviations from (e) the expected Nambu-Goto-like behavior, and, for the first time, evidence for (e) contributions that can be described by a term proportional to the extrinsic curvature of the effective string worldsheet. Such term is allowed by (e) Lorentz invariance and its presence in the infrared regime of the U (1) model was indeed predicted by (e) Polyakov several years ago. Results show that this term scales as expected according to (e) Polyakov's solution, and becomes (=) the dominant contribution to (e) the effective string action in the continuum limit. Authors also demonstrate analytically that the corrections to (e) the confining potential induced by (e) the extrinsic curvature term can be related to (e&eb) the partition function of (e) the massive perturbation of (e) a c=1 bosonic conformal field theory. The implications of results for SU (N) Yang-Mills theories in (e&eb) three and in four spacetime dimensions are discussed. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th) Journal reference: JHEP 1501 (2015) 105 DOI: 10.1007/JHEP01 (2015)105 Report number: IFT-UAM/CSIC-14-052 Cite as: arXiv:**

**1406.5127 [hep-lat] (or arXiv: 1406.5127v2 [hep-lat] for this version) A different kind of string Michele Caselle, Marco Panero, Roberto Pellegrini, Davide Vadacchino STOP In the past ten years, the ideas of supersymmetry have been profitably applied to (e&eb) many nonrelativistic quantum mechanical problems. In particular, there is now a much deeper understanding of why certain potentials are analytically solvable and (e&eb) an array of powerful new approximation methods for (e) handling potentials which are not (e) exactly solvable. In this report, authors review the theoretical formulation of (e) supersymmetric quantum mechanics and discuss many applications. Exactly solvable potentials can be understood in terms of (e&eb) a few basic ideas which include (e) supersymmetric partner potentials (e&eb) shape invariance and (e&eb) operator transformations. Familiar solvable potentials all have (e) the property of shape invariance. Authors describe new exactly solvable shape invariant potentials which include (e) the recently discovered self-similar potentials as (=) a special case. The connection between inverse scattering,(e&eb) isospectral potentials and (e&eb) supersymmetric quantum mechanics is discussed and multisoliton solutions of (e) the K dV equation are constructed. Approximation methods are also discussed within (eb) the framework of supersymmetric quantum mechanics and in particular it is shown that a supersymmetry inspired WKB approximation is (=) exact for a class of shape invariant potentials. Supersymmetry ideas give (eb) particularly nice results for the tunneling rate in (eb) a double well potential and for improving (eb+) large N expansions. Authors also discuss the problem of (e) a charged Dirac particle in (eb) an external magnetic field and other potentials in terms of (e&eb) supersymmetric quantum mechanics. Finally, authors discuss structures more general than supersymmetric quantum mechanics such as(=) parasupersymmetric quantum mechanics in which there is symmetry between a boson and (e&eb) a para-fermion of order p. Copyright © 1995 Published by Elsevier B.V. Physics Reports Volume 251, Issues 5–6, January 1995, Pages 267–385Supersymmetry and quantum mechanics Fred Coopera, Avinash Khareb, Uday Sukhatmec STOP The free energy of (e) the maximally**

**supersymmetric SU (N) gauge theory at temperature T is expected to scale, in the large-N limit, as (=) N2T4 times a function of the 't Hooft coupling, f (gYM2N). In the strong coupling limit the free energy has been deduced from (e) the near-extremal 3-brane geometry, and its normalization has turned out to be Full-size image (<1 K) times that found in (eb) the weak coupling limit. In this paper authors calculate the leading correction to (e) this result in inverse powers of (e) the coupling, which originates from (e) the R4 terms in the tree-level effective action of (e) type IIB string theory. The correction to Full-size image (<1 K) is positive and of (e) order Full-size image (<1 K). Thus, f (gYM2N) increases as (e&eb) the 't Hooft coupling is decreased, in accordance with (e&eb) the expectation that is should be approaching 1 in (eb) the weak coupling limit. Authors also discuss similar corrections for (e) other conformal theories describing (eb) coincident branes. In particular, authors suggest (eb) that the coupling independence of (e) the near extremal entropy for (e) D1-branes bound to (e) D5-branes is related to (e&eb) the vanishing of (e) the Weyl tensor of AdS3 × S3. Copyright © 1998 Published by Elsevier B.V. Nuclear Physics B Volume 534, Issues 1–2, 23 November 1998, Pages 202–222Coupling constant dependence in the thermodynamics of N = 4 supersymmetric Yang-Mills theory Steven S. Gubsera, Igor R. Klebanovb, 1, Arkady A. Tseytlinc, 2 STOP A new approach to (e) N=2 supersymmetry based on (e) the concept of harmonic superspace is proposed and is used to give (eb) an unconstrained superfield geometric description of N=2 super Yang-Mills and supergravity theories as (=) well as of matter N=2 hypermultiplets. The harmonic N=2 superspace has (e) an independent coordinates, in addition to the usual ones, the isospinor harmonics ui+or- on (e&eb) the sphere SU (2)/U (1). The role of ui+or- is to relate the SU (2) group realised on the component fields to (e&eb) a U (1) group acting on (e&eb) the relevant superfields. Their introduction makes it possible to SU (2)-covariantise (eb (t), e&eb) the notion of Grassmann analyticity. Crucial for the construction is the existence of (e) an analytic subspace of the general harmonic N=2 superspace. The hypermultiplet superfields and (e&eb) the true prepotentials (pre-**

**prepotentials) of N=2 super Yang-Mills and supergravity are (=) unconstrained superfunctions over (e&eb) this analytic subspace. The pre-prepotentials have (e) a clear geometric interpretation as gauge connections with (e&eb) respect to the internal SU (2)/U (1) directions. A radically new feature arises: the number of gauge and (e&eb) auxiliary degrees of freedom becomes (=) infinite while (e) the number of physical degrees of freedom remains (=) finite. Other new results are the massive N=2 Yang-Mills theory and (e&eb) various off-shell self-interactions of hypermultiplets. The propagators for matter and Yang-Mills superfields are given. Unconstrained N=2 matter, Yang-Mills and supergravity theories in harmonic superspace A Galperin, E Ivanov, S Kalitzin, V Ogievetsky and E Sokatchev STOP By assuming the existence of a novel multipronged string state for (e) D-particles interacting with (e&eb) D-brane intersections in type IIA string theory, authors are able to derive a quantum mechanical description of (e) supersymmetric Reissner-Nordstrom black holes. A supersymmetric index calculation provides (eb) evidence for (e) this conjecture. The quantum mechanical system becomes (=) two decoupled conformal quantum mechanical systems in (eb) the low energy limit. The conformal quantum mechanics has expected properties of (e) a dual description of string theory on (e&eb) AdS2×S2. Subjects: High Energy Physics - Theory (hep-th) Journal reference: JHEP 0001:006, 2000 DOI: 10.1088/1126-6708/2000/01/006 Report number: LBNL-43734 Cite as: arXiv:hep-th/9907101 (or arXiv:hep-th/9907101v5 for this version) Supersymmetric Quantum Mechanical Description of Four Dimensional Black Holes Julie D. BlumSTOP Using recently proposed soliton equations authors derive a basic identity for the scaling violation of (e) N=2 supersymmetric gauge theories Σiai∂F/∂ai−2F=8πib1u. Here F is (=) the prepotential, ai's are (=) the expectation values of (e) the scalar fields in (eb) the vector multiplet, u=1/2 Tr <ϕ2> and b1 is the coefficient of (e) the one-loop β-function. This equation holds in (eb) the Coulomb branch of all N=2 supersymmetric gauge theories coupled with (e&eb) massless matter. TOHRU EGUCHI and SUNG-KIL YANG, Mod Phys. Lett. A 11, 131 (1996).**

**DOI: 10.1142/S0217732396000151 PREPOTENTIALS OF N=2 SUPERSYMMETRIC GAUGE THEORIES AND SOLITON EQUATIONS STOP The Coulomb branch of N = 2 supersymmetric gauge theories in four dimensions is described in general by (e) an integrable Hamiltonian system in (eb) the holomorphic sense. A natural construction of such systems comes from (e) two-dimensional gauge theory and (e&eb) spectral curves. Starting from this point of view, authors propose an integrable system relevant to (e&eb) the N = 2 SU (n) gauge theory with (e&eb) a hypermultiplet in the adjoint representation, and offer (eb) much evidence that it is correct. The model has (e) a Full-size image (<1 K) S-duality group (with the central element −1 of Full-size image (<1 K) acting as (=) charge conjugation); Full-size image (<1 K) permutes the Higgs, confining, and oblique confining phases in the expected fashion. We also study more exotic phases. Copyright © 1996 Published by Elsevier B.V. Nuclear Physics B Volume 460, Issue 2, 5 February 1996, Pages 299–334Supersymmetric Yang-Mills theory and integrable systems Ron Donagia, 1, Edward Wittenb, 2 STOP A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 present a first step towards generalizing (eb) the work of Seiberg and Witten on N = 2 supersymmetric Yang-Mills theory to (e&eb) arbitrary gauge groups. Specifically, authors propose a particular sequence of (e) hyperelliptic genus n − 1 Riemann surfaces to underly (e&eb) the quantum moduli space of SU (n)N = 2 supersymmetric gauge theory. These curves have (e) an obvious generalization to (e) arbitrary simply laced gauge groups, which involves (e&eb) the A-D-E type simple singularities. To support proposal, authors argue that the monodromy in (e&eb) the semiclassical regime is (=) correctly reproduced. We also give some remarks on a possible relation to string theory. Copyright © 1995 Published by Elsevier B.V. Physics Letters B Volume 344, Issues 1–4, 26 January 1995, Pages 169–175Simple singularities and N = 2 supersymmetric Yang-Mills theory A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 STOP Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 review and extend earlier work that uses (e) the AdS/CFT correspondence to relate the**

**black-hole–black-string transition of gravitational theories on a circle to (e&eb) a phase transition in maximally supersymmetric (1 + 1)-dimensional SU(N) gauge theories at (eb) large N, again compactified on (e&eb) a circle. Authors perform gravity calculations to determine (eb) a likely phase diagram for (e) the strongly coupled gauge theory. Authors then directly study the phase structure of (e) the same gauge theory, now at (eb) weak't Hooft coupling. In the interesting temperature regime for (e) the phase transition, the (1 + 1)-dimensional theory reduces to (e&eb) a (0 + 1)-dimensional bosonic theory, which they solve using (e) Monte Carlo methods. Authors find strong evidence that the weakly coupled gauge theory also exhibits (eb) a black hole–black string-like phase transition in (eb) the large N limit. Authors demonstrate that a simple Landau–Ginzburg-like model describes (eb) the behaviour near the phase transition remarkably well. The weak coupling transition appears to be (=) close to the cusp between a first-order and a second-order transition. Black-hole–black-string phase transitions in thermal (1 + 1)-dimensional supersymmetric Yang–Mills theory on a circle Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 Published 29 October 2004 • 2004 IOP Publishing Ltd Classical and Quantum Gravity, Volume 21, Number 22STOP Monitoring the orbits of stars around (e&eb) Sgr A* offers(eb)the possibility of detecting (e&eb) the precession of (e) their orbital planes due to (e) frame dragging, of (e) measuring the spin and quadrupole moment of (e) the black hole, and of (e)testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and (e&eb) tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are (=), in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for which an optical interferometer, such as GRAVITY, will detect orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to (e&eb) the timescale of precession due to (e) the quadrupole moment of (e) the black hole. As a result, this non-conservative effect is (=) a**

**potential source of systematic uncertainty in (eb) testing the no-hair theorem with (e&eb) stellar orbits. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI:10.1088/0004-637X/777/1/57 Cite as:arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this version) Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP Authors perform a phase space analysis of (e) a generalized modified gravity theory with (e&eb) non-minimal coupling between geometry and (e&eb) matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that (eb) in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, Authors prove that there is (=) an attractor solution to describe (eb) the late time accelerating universe when (e) the modified gravity is chosen in (eb) a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and supernova type Ia data sets, Authors find that (eb) the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as:arXiv: 1512.09281 [gr-qc] (or arXiv: 1512.09281v1 [gr-qc] for this version) Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang study the P−V criticality and phase transition in (eb) the extended phase space of (e) charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of the system and (e&eb) its conjugate quantity is (=) the thermodynamic volume of (e) the black hole. The black**

**holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb)no P−V criticality and phase transition appear, while for the black holes with a spherical horizon, even when (e) the charge of the black hole is absent, the P−V criticality and the small black hole/large black hole phase transition (eb) will appear, but it happens only in (eb) d=5 dimensions; when (e) the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to (e) have the P−V criticality and the small black hole/large black hole phase transition, there exists (eb) an upper bound for (e) the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the Gauss-Bonnet coefficient and Q is the charge of the black hole. Authors calculate the critical exponents at the critical point and find that for all cases, they are (=) the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09 (2013)005 Cite as: arXiv: 1306.6233 [gr-qc] (or arXiv: 1306.6233v4 [gr-qc] for this version) P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP paper studies the whole process of (e) gravitational collapse and (e&eb) accretion of (e) a massless scalar field in (eb) asymptotically flat spacetime. Two kinds of initial configurations are considered. One is the initial data without (e) black hole; the other contains (e) a black hole. Under suitable initial conditions, Authors find that multi-horizon will appear which mean (eb) that the initial black hole formed by (e) gravitational collapse or existing at the beginning will instantly expand and suddenly grow (e&eb, eb), rather than grows gradually in the accretion process. A new type of critical behavior is found around (e&eb) the instant expansion. The numerical computation shows (eb) that the critical exponents are (=) universal. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Computational Physics (physics. Comp-ph) Cite as: arXiv: 1604.03363 [gr-qc] (or arXiv: 1604.03363v1 [gr-qc] for this version)**

**Multi-horizon and Critical Behavior in Gravitational Collapse of Massless Scalar Zhoujian Cao, Rong-Gen CAI, Run-Qiu Yang STOP In this paper, Authors present the results on the gravitational collapse of (e) charged massless scalar field in asymptotically flat spacetime with(e&eb) a perfectly reflecting wall. Differing from previous works, Authors study the system in (eb) the double null coordinates, by which we could simulate (e&eb) the system until the black hole forms with (e&eb) higher precision but less performance time. Authors investigate the influence of charge on (e&eb) the black hole formation and the scaling behavior near the critical collapses. The gapless and gapped critical behaviors for (e) black hole mass and charge are studied numerically. Authors find that they satisfy (eb) the scaling laws for critical gravitational collapse but the gapped critical behavior is different from (e) its AdS counterpart. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1602.00112 [gr-qc] (or arXiv: 1602.00112v1 [gr-qc] for this version) Multiple critical gravitational collapse of charged scalar with reflecting wall Rong-Gen CAI, Run-Qiu Yang STOP In this paper Authors investigate the (2+1)-dimensional topological non-Fermi liquid in (eb) strongly correlated electron system, which has (e) a holographic dual description by (e) Einstein gravity in (3+1)-dimensional anti-de Sitter (AdS) space-time. In a dyonic Reissner-Nordstrom black hole background, Authors consider a Dirac fermion coupled to (e&eb) the background U (1) gauge theory and an (e&eb) intrinsic chiral gauge field bM induced by (e) chiral anomaly. UV retarded Green's function of (e) the charged fermion in (eb) the UV boundary from (e) AdS4 gravity is calculated, by imposing (e&eb) in-falling wave condition at the horizon. Authors also obtain IR correlation function of (e) the charged fermion at (eb) the IR boundary arising from (e) the near horizon geometry of (e) the topological black hole with (e&eb) index k=0,±1. By using (e) the UV retarded Green's function and IR correlation function Authors analyze the low frequency behavior of (e) the topological non-Fermi liquid at zero and (e&eb) finite temperatures, especially the relevant non-Fermi liquid behavior near (eb) the quantum critical region. In addition, Authors**

**find that when k=±1, the effective mass of (e) bulk Dirac fermion at IR fixed point is (=) topologically quantized, consequently the Fermi momentum presents (eb) a discrete sparse distributed pattern, in contrary to (e&eb) the Ricci flat case with k=0. As demonstrated examples, Authors calculate the spectral functions of (e) topological non-Fermi liquid due to (e) anomalous Hall effect (AHE) (|b→|≠0,b 0=0) and chiral magnetic effect (CME) (|b→|=0,b0≠0), which can be tested in topological materials such as (=) topological insulator and (e&eb) topological metal with (e&eb) strong couplings. Subjects: High Energy Physics - Theory (hep-th); Strongly Correlated Electrons (cond-mat.str-el); General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv: 1601.03865 [hep-th] (or arXiv: 1601.03865v1 [hep-th] for this version) Topological Non-Fermi Liquid Rong-Gen Cai, Yong-Hui Qi, Yue-Liang Wu, Yun-Long Zhang STOP The unusual magnetic materials are significant in both science and technology. However, because of (e) the strongly correlated effects, it is difficult to understand their novel properties from (e) theoretical aspects. Holographic duality offers (eb) a new approach to understanding such systems from (e) gravity side. This paper will give a brief review of Authors recent works on the applications of holographic duality in understanding (e&eb) unusual magnetic materials. Some quantitative compare between holographic results and (e&eb) experimental data will be shown and some predictions from (e) holographic duality models will be discussed. Comments: For Second LeCosPA Symposium at Taiwan National University, Taipai, Dec. 14-18, 2015 Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1601.02936 [hep-th] (or arXiv:1601.02936v1 [hep-th] for this version) Understanding strongly coupling magnetism from holographic duality Rong-Gen Cai, Run-Qiu Yang STOP It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a**

**Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3**

**[astro-ph.CO] for this version) Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang**

# Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

## Module Number Six

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-465) 1990. We demonstrate its applicability cosmology/ quantum computation economics etal. For the first time, physicists couple distant nuclear spins using a single electron July 11, 2016. Some sentences have been deleted due to spatial restrictions. Please excuse me. For the first time, researchers at the University of Basel in Switzerland have coupled the nuclear spins of distant atoms using (e) just a single electron. Three research groups took part in this complex experiment, the results of which have now been published in the journal Nature Nanotechnology. In most materials, the nuclear spins of (e) neighboring atoms have (e) only a very weak effect on one another, as (=) the tiny nuclei are located deep within (e&eb) the atoms. This is not the case in metals, however, where (e) some of the electrons are able to move freely. The electron spins are able to couple nuclear spins at (eb) relatively large distances from one another. Named after four physicists, this RKKY interaction was discovered in the 1950s. Nuclear spins linked by (e&eb) an individual electron spin Now, an experiment by researchers from the Department of Physics at the University of Basel has for the first time succeeded in demonstrating this mechanism with (e&eb) a single electron, describing it using (e) quantum theory. The team led by Prof. Richard Warburton trapped a single electron inside (eb) a quantum dot. With use of a method developed in Basel to measure (eb) nuclear spin resonance, they showed that the electron coupled (e&eb) the spins of nuclei up to five nanometers apart - a huge distance in the world of (e) nuclear spins. The results are particularly relevant to (e&eb) the development of spin qubits: these seek to use (e) electron spins to carry (e&eb) information, but the interaction with (e&eb) the nuclei limits (e) the stability of the quantum information. Concentrated blast of physics More information: Role of the electron spin in determining the coherence of the nuclear spins in a quantum dot, Nature Nanotechnology, DOI: 10.1038/nnano.2016.114 Journal reference: Nature Nanotechnology search and more info website Provided by: University of Basel Read more at: http://phys.org/news/2016-07-physicists-couple-distant-nuclear-electron.html#jCp A new spin on quantum computing: Scientists train electrons with microwaves February 15, 2016 A silicon sample doped with (e&eb) bismuth atoms (left image) that is (=) just 150 nanometers thick is fitted with (e&eb) a superconducting resonator that includes (e) a capacitor (black, in left image; light gray in center image) and an inductive wire (red line … In what may provide (eb) a potential path to processing information in (eb) a quantum computer, researchers have switched an intrinsic property of electrons from an excited state to (e&eb) a relaxed state on demand using (e) a device that served as (=) a microwave "tuning fork." The team's findings could also lead to (eb) enhancements in magnetic resonance techniques, which are (=) widely used to explore (e) the structure of materials and biomolecules, and for medical imaging. The**

**international research team, which included scientists at the Department of Energy's Lawrence Berkeley National Laboratory (Berkeley Lab), demonstrated how to dramatically increase the coupling of (e) microwaves in a specially designed (e&eb) superconducting cavity to (e) a fundamental electron property called spin—which, like a coin, can be flipped. By zapping an exotic silicon material developed at Berkeley Lab with the microwaves, they found that they could rapidly change (e&eb) the electron spins from (e) an excited state to a relaxed, ground state by (e) causing (eb) the electrons to emit (eb) some of their energy in the form of (e) microwave particles known as photons. Left on their own, the electron spins would be extremely unlikely to flip back to (e) a relaxed state and to also emit (eb) a photon - the natural rate for this light-emitting effect, known as the Purcell effect, is about once every 10,000 years. The experiment demonstrated an accelerated, controllable relaxation of (e) electron spins and the release of (e) a microwave photon in about 1 second, said Thomas Schenkel, a physicist in Berkeley Lab's Accelerator Technology and Applied Physics Division who led the design and development of the silicon-bismuth sample used in the experiment. A new spin on quantum computing: Scientists train electrons with microwaves Like the classical tuning fork slung from guitar strings in this photo, scientists used (e) a device that worked like a microwave tuning fork to switch (e&eb) a fundamental property of electrons on (e&eb) demand in a silicon-bismuth sample developed at Berkeley Lab. Credit: Flickr/Pierre Guinoiseau "It's like a juggler who throws the balls up, and the balls come down 1,000 times faster than normal, and they also emit (eb) a microwave flash as they drop," he said. The results were published online Feb. 15 in the journal Nature. "Our results are highly significant for (e&eb) quantum information processing," said Patrice Bertet, a quantum electronics scientist at the French Atomic Energy Commission (CEA) who led the experiment. "Indeed, they are a first step toward (e&eb) the strong coupling of (e) individual electron spins to microwave photons, which could form (eb) the basis of new spin-based quantum computer architecture." John Morton, a professor at the London Center for Nanotechnology and co-author of the study, said, "Our ultimate aim is to find a link between quantum information that is fixed and (e&eb) quantum information that can be transported by photons." In today's computers, information is stored as (=) individual bits, and each bit can either be (=) a one or a zero. Quantum computers, though, could conceivably be (=) exponentially more powerful than modern computers because (e) they would use a different kind of bit, called a qubit, which because of (e) the weird ways of quantum mechanics can simultaneously behave (e&eb) as both a one and a zero. A coupled array of qubits would allow (eb) a quantum computer to perform (e&eb) many, many calculations at (eb) the same time, and electron spins are (=) candidates for qubits in (eb) a quantum computer. The latest study shows how the microwave photons could work in concert with (e&eb) the spins of electrons to move information in a new type of computer. "What we need now is ways to wire up these systems—to couple (e&eb) these spins together," Morton said. "We need to make coupled qubits that can perform (e&eb) computations." In the experiment, conducted at CEA in France, a small sample of a highly purified form of (e) silicon was implanted with (e&eb) a matrix of bismuth atoms, and a superconducting aluminum circuit was deposited on (e&eb) top to create (eb) a high-quality resonant cavity that allowed (eb) precise tuning of the microwaves. The electron spins of the bismuth atoms were then flipped into (e&eb) the excited, "spin-up" state. The microwave cavity was then tuned (e&eb) like a musical tuning fork, to (e) a particular resonance that coaxed the spins into emitting (eb) a photon as they flipped back to (e) a relaxed state. The cavity boosted the number of states into (e&eb) which a photon can be (=) emitted, which greatly increased (eb) the decay rate for (e) the electron spins in (eb) a controllable way. The technique is much like buying more lottery tickets to increase (eb) your chances of winning, Morton said. A new spin on quantum computing: Scientists train electrons with microwaves From left are, Berkeley Lab scientists Thomas Schenkel, Qing Ji and Peter Seidl at the NDCX-II (Neutralized Drift Compression Experiment II), which produces (eb) powerful ion beams. Researchers are exploring how to use NDCX-II to process exotic … The large bismuth atoms embedded in (eb) the silicon sample provided (eb) the electrons with (e&eb) unique spin properties that enabled (eb) the experiment. Schenkel said that implanting the bismuth atoms into (e&eb) the delicate silicon framework, a**

**process known as "doping," was "like squeezing bowling balls into (e&eb) a lattice of ping-pong balls." "We did a new trick with silicon. People wouldn't expect you could squeeze (e) anything new out of (e) silicon," Schenkel said. "Now we're looking into further improving (eb) bismuth-doped silicon and into (e&eb) tailoring the spin properties of (e) other materials, and using (e) this experimental technique for these materials." To enhance the performance of materials used (e) in future experiments, Schenkel said it will be necessary to improve (eb) the doping process so it is less damaging to (e) the silicon lattice. Also, the implantation process could be designed to produce (eb) regularly spaced arrays of (e) individual electron spins that would be more useful for (e) quantum computing than a concentrated ensemble of (e) electron spins. "We are now doing experiments on processing (e&eb) this and other materials at higher temperature and pressure with nanosecond ion pulses at NDCX-II, one of the accelerators here at Berkeley Lab," Schenkel said. "There are indications that it will improve (eb) the overall spin quality." Researchers said the latest research could potentially prove useful in boosting (eb) the sensitivity of scientific techniques like nuclear magnetic resonance spectroscopy and dynamic nuclear polarization, useful for a range of experiments, and could also shorten (e) experimental times by manipulating (e&eb) spin properties. "You need a way to reset (e&eb) spins—the ability to cause (eb) them to relax on demand to improve (eb) the rate at which you can repeat an experiment," Morton said. Bertet said it may be possible to further accelerate (eb) the electron-flipping behavior to (e) below 1 millisecond, compared to (e) the 1-second rate in the latest results. "This will then open the way to many new applications," he said. Explore further: Bismuth provides perfect dance partners for quantum computing qubits More information: A. Bienfait et al. Controlling spin relaxation with a cavity, Nature (2016). DOI: 10.1038/nature16944 Nature | Letter Controlling spin relaxation with a cavity A. Bienfait, J. J. Pla, Y. Kubo,X. Zhou, M. Stern, C. C. Lo, C. D. Weis, T. Schenkel, D. Vion,D. Esteve, J. J. L. Morton& P. Berte 531, 74–77 (03 March 2016) doi:10.1038/nature16944 Spontaneous emission of radiation is (=) one of the fundamental mechanisms by which an excited quantum system returns to (e) equilibrium. For spins, however, spontaneous emission is (=) generally negligible compared to (e&eb) other non-radiative relaxation processes because of (e) the weak coupling between the magnetic dipole and (e&eb) the electromagnetic field. In 1946, Purcell realized that the rate of spontaneous emission can be greatly enhanced by (e) placing the quantum system in a resonant cavity. This effect has since been used extensively to control (e&eb) the lifetime of atoms and semiconducting heterostructures coupled to (e&eb) microwave2 or optical3, 4 cavities, and is essential for (e) the realization of high-efficiency single-photon sources5. Here authors report the application of this idea to (e&eb) spins in solids. By coupling donor spins in silicon to (e&eb) a superconducting microwave cavity with (e&eb) a high quality factor and a small mode volume, authors reach the regime in which spontaneous emission constitutes (e) the dominant mechanism of spin relaxation. The relaxation rate is increased by (e) three orders of magnitude as the spins are tuned to (e) the cavity resonance, demonstrating that energy relaxation can be controlled (e&eb) on demand. Results provide a general way to initialize (eb) spin systems into their ground state and therefore have applications in (eb) magnetic resonance and (e&eb) quantum information processing6. They also demonstrate that the coupling between the magnetic dipole of a spin and (e&eb) the electromagnetic field can be enhanced up to (e) the point at which quantum fluctuations have (e) a marked effect on (e&eb) the spin dynamics; as such, they represent (eb) an important step towards the coherent magnetic coupling of (e) individual spins to (e) microwave photons Scientists obtain evidence of many-body localization in a closed quantum system July 8, 2016 by Olivia Meyer-Streng Science Cover Vol 352, Issue 6293, 24 June 2016, to highlight the research published in this issue (see original publication). The cover shows an artist's view of the experiment, illustrating the persistence of the density step of (e) an atomic distribution in an optical lattice. During equilibration ordinary many-body systems lose (e) all information about (e&eb) the initial state. Every morning authors experience an example for (e) this behaviour. Milk poured into a cup of coffee mixes (e&eb) perfectly and after some time it is impossible to say how exactly the two fluids were (=) put together. The same behaviour holds for (e) almost all quantum systems. However,**

**recently a new phenomenon called "many-body localization" has been predicted theoretically, which allows (eb) well insulated quantum systems to (e) preserve memory of the initial state forever. Now a team of scientists around Dr. Christian Groß and Professor Immanuel Bloch (Director at MPQ and Chair of Quantum Optics at LMU Munich), in cooperation with David Huse (Princeton University), has obtained evidence of such a behaviour in a two-dimensional quantum system of cold rubidium atoms trapped in (e&eb) an optical lattice. The scientists observed that – beyond a certain degree of disorder imprinted on (eb) the particle ensemble in the beginning – the system would relax into (e&eb) a steady state still containing (e) detailed microscopic information about its past. "We were able to observe the transition from a thermalized state into (e&eb) a many-body localized phase", Christian Groß points out. "It is the first observation of that kind in a regime that is not accessible with (e&eb) state-of-the-art simulations on classical computers." The experiment is not only of fundamental interest; the results might also lead to (eb) new ways for storing quantum information. Motivated by the foundational problem of how interacting particles behave (e&eb) in a disordered system, in the 1950s the American physicist Philip Warren Anderson discovered the famous localization phenomenon for (e) non-interacting particles, now called "Anderson localization". Here, disorder prevents (e) the particles to move and consequently all transport is (=) stopped. But what happens when disorder comes together with (e&eb) interactions? Will interactions lead to (eb) transport and thermalization, or will the localization persist even at (eb) high energies? So far, there is no theoretical model that faithfully predicts (eb) the evolution of a closed quantum system in more than one dimension under (e&eb) these conditions, although, the possibility for (e)localization has been theoretically suggested. In order to investigate these questions experimentally, strict requirements on controllability and perfect isolation of (e) the system have to be fulfilled. In the experiment described here, ultracold rubidium atoms are loaded into (eb) an optical lattice, a microscopic array of light traps formed by interfering (e&eb) laser beams. The disorder is created by projecting (e&eb) a computer-generated random light pattern onto (e&eb) the optical lattice. As a result of this "disorder", the depth of each microscopic trap in the lattice potential varies from site to (e) site. In fact, Prof. Bloch's group has advanced the experimental tools to such perfection that they can steer the position of (e) the atoms in the artificial light crystal and (e&eb) the interaction between (e&eb) them almost at will. With a high resolution microscope that detects the emitted fluorescence light, the position of each atom can be observed with (e&eb) high precision. Additionally, the atomic density distribution of the initial state can be (=) controlled, and for variable time intervals the evolution of the distribution can be (=) measured with high precision. These tools at hand, non-thermalizing behaviour can be probed in (eb) a conception ally very simple way. Any thermalized state of a closed system reflects (e&eb) the symmetry of its container – like, for example, water spilled into (e&eb) a round bowl would immediately cover (e) the whole bottom. In analogy, the scientists generate (eb) a density step in the initially prepared sample by (e) blowing away half of the atomic distribution with laser radiation. Then they watch how the remaining particles migrate into (e&eb) the empty half. For small disorders, the initially prepared density step is smeared out (e) fast and the initially empty and filled halves become (=) indistinguishable. However, when the measurement is repeated for (e) strong disorder, traces of the initial state remain and the system does not (e) relax to a thermal state even for (e) very long times. "We observe a fairly sharp onset of non-thermalizing behaviour above (e) a critical value", says Christian Groß. "This absence of thermalization is remarkable because (e) thermalization persists in a system of interacting particles, even at (eb) the high energies probed in the experiment." The scientists interpret their observation as (=) the onset of many-body localization in (eb) the atomic system. This is of fundamental interest because it means (eb) the breakdown of equilibrium statistical mechanics. On the other hand, the persistence of initial state information could be used as (=) a source for quantum information technologies. "It should also be emphasized that we obtain these results for a system size that is far beyond (e) numerically accessible scales", says Jae-yoon Choi, postdoc at the experiment. More information: J.-y. Choi et al. exploring the many-body localization transition in two dimensions, Science (2016). DOI: 10.1126/science.aaf8834 Journal reference: Science search and more**

**info website Provided by: Max Planck Society search and more info website Bosons refusing to thermalize in 2D Messy, interacting quantum-mechanical systems are difficult to analyze theoretically. In a single spatial dimension, the calculations are (=) still tractable, and experiments have recently confirmed (eb) the prediction that sufficiently strong disorder can disrupt (e&eb) the transport of interacting particles. In two dimensions, however, the theoretical blueprint is missing. Choi et al. used (e) single-site imaging of cold 87Rb atoms in (eb) an optical lattice to show (eb) that similar localization occurs in (eb) two-dimensional (2D) systems. The study highlights the power of quantum simulation to solve problems that are currently inaccessible to classical computing techniques. Science, this issue p. 1547 Abstract A fundamental assumption in statistical physics is that generic closed quantum many-body systems thermalize under (e&eb) their own dynamics. Recently, the emergence of many-body localized systems has questioned (e) this concept and challenged our understanding of the connection between statistical physics and (e&eb) quantum mechanics. Here authors report on the observation of a many-body localization transition between thermal and (e&eb) localized phases for (e) bosons in a two-dimensional disordered optical lattice. With our single-site–resolved measurements, authors track the relaxation dynamics of (e) an initially prepared out-of-equilibrium density pattern and find strong evidence for (e) a diverging length scale when (e) approaching the localization transition. Experiments represent a demonstration and in-depth characterization of many-body localization in (eb) a regime not accessible with (e&eb) state-of-the-art simulations on classical computers. Pelican instrument provides crucial experimental evidence of unusual quantum state July 8, 2016 Artist's illustration. Credit: Australian Nuclear Science and Technology Organisation (ANSTO) Experiments using inelastic neutron scattering at the Australian Centre for Neutron Scattering have found indications (eb) of a possible new quantum spin state in a novel antiferromagnetic material barium ytterbium zinc oxide (Ba3Yb2Zn5O11) which provides (eb) both a challenge and validation of (e) the third law of thermodynamics. The quantum mechanical information gained from (e) the experiment is crucial not only in understanding the quantum state of the specific material but also provides (eb) better understanding of high temperature superconductivity and ultimately offers (eb) guidance in the design of materials for (e) quantum computing. A research team led by Associate Professor Takatsugu Masuda at the Institute for Solid State Physics at the University of Tokyo with collaborators from ANSTO, Osaka University and Taiwan's National Cheng Kung University, have performed a successful experiment using Pelican, time-of-flight cold neutron spectrometer at ANSTO. The results have just been published in Physical Review B as a rapid Communication and an Editors' Suggestion. "The great success of this challenging experiment relied on (e) the excellent performance of the Pelican instrument, wonderful technical support at ANSTO and strong expertise in the relevant field from A/Prof Masuda's group," said Dr Dehong Yu and Dr Richard Mole, co-authors and Pelican instrument scientists. Pelican instrument provides crucial experimental evidence of unusual quantum state INS spectra measured at 1.5 K (a), 6 K (b), 12 K (c), and 40 K (d) using setup I. (e) INS spectrum measured at 1.5 K using the setup II. (f) Q dependence of the integrated intensity obtained from spectrum in (e). Red solid curve is the …….(deleted due to spatial constraints) Barium ytterbium zinc oxide (Ba3Yb2Zn5O11) is a novel antiferromagnetic model system having (e) a "breathing pyrochlore lattice" consisting of (e) arrays of alternating large and (e&eb) small corner-sharing tetrahedra of ytterbium atoms. Earlier research had suggested (eb) the material was (=) an isotropic spin half (S= ½) system. However, four low energy excitations measured on the sample with the Pelican instrument at a temperature of 1.5 K clearly disagree with (e) this designation and the associated theoretical model based on (e) the isotropic spins having simple interactions. "These observations suggested (eb) a much more complex spin state in the system," said Yu. "The experiment is interesting and important because as you start going very low in temperature, quantum mechanics takes over (e) system rather than thermodynamics, and you observe unusual properties," said Mole. "According to low temperature heat capacity measurements, the material does not (e) have zero entropy or a single ground state until you reach 63 milli Kelvin or -272.94 Celsius, before it starts obeying the third law." Pelican instrument provides crucial experimental evidence of (e) unusual quantum state**

**Breathing pyrochlore lattice formed by (e) A3+ ions as determined by Kimura, Nakatsuju and Kimura in previous work The third law of thermodynamics states that for pure crystals as the temperature of a system approaches absolute zero K, (eb) its entropy or disorder approaches zero. "The research clearly indicates (eb) that below this temperature the material may be in (eb) an exotic state, such as (=) a liquid-like spin state, in which the arrangement of spins behaves (e&eb) like a liquid. In this exotic state, novel phenomena may show up and be detectable; for example, charge and spin separation of electrons may be (=) possible," said Yu. The researchers suggest that further investigations, experimentally and theoretically, are required to ultimately determine the true identity of (e) the ground state of the material. "Discovering the existence of regular tetrahedral spin made it well worth coming all the way to Australia to make use of this neutron experimental apparatus," said Associate Professor Masuda. Many Japanese researchers are using ANSTO's neutron scattering instruments through the peer review system, while their nuclear facilities are offline. The outcome of this particular experiment is also part of the collaboration agreement between ANSTO and the University of Tokyo. More information: T. Haku et al. Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet, Physical Review B (2016). DOI: 10.1103/PhysRevB.93.220407 Journal reference: Physical Review B search and more info website: provided by: Australian Nuclear Science and Technology Organisation (ANSTO) search and more info website Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 T. Haku, K. Kimura, Y. Matsumoto, M. Soda, M. Sera, D. Yu, R. A. Mole, T. Takeuchi, S. Nakatsuji, Y. Kono, T. Sakakibara, L.-J. Chang, and T. Masuda Phys. Rev. B 93, 220407(R) – Published 24 June 2016 Authors study low-energy excitations in (eb) the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 by (e) a combination of inelastic neutron scattering (INS) and (e&eb) thermodynamical property measurements. The INS spectra are quantitatively explained by (e) spin-12 single-tetrahedron model having (e) XXZ anisotropy and Dzyaloshinskii-Moriya interactions. This model has a twofold degeneracy of (e) the lowest-energy state per tetrahedron and well reproduces (eb) the magnetization curve at (eb) 0.5 K and heat capacity above 1.5 K. At lower temperatures, however, authors observe a broad maximum in (eb) the heat capacity around 63 mK, demonstrating (eb) that a unique quantum ground state is selected due to (e) extra perturbations with (e&eb) an energy scale smaller than the instrumental resolution of INS. Revised 13 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.220407 © 2016 American Physical Society Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA (Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor. In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface. As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone. Such planets are all the "right" distance from (e) their star, and hopefully, also have water. But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone. For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement. But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway. But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain. It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of**

**reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life. Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface. The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113 Abstract The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question. Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry. This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry. Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments. Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy. The electronic and (e&eb) structural variability among them is extraordinary. The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations. The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface. The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers. Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp**

## 57. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_6}{dt} = -b_6 Q_6 + g_6(g_6)Q_6$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_6(0) = 0, \quad g_6(P_6) > 0, \ P_6 \ 0, \qquad \lim_{P_6 \to \infty} g_6(P_6) = (K_6)_6 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_6 + (\rho_a)_6 + D_6(a_6, P_6, Q_6)\rho_6 + 0, \qquad 0 < a_6 \ \infty,$$

$$\rho_6(t, 0) = \int_0^\infty n_6(a_6, P_6, Q_6)\rho_6(t_6, a_6)da_6$$

Where $D_6$ is the mortality function and $n_6$ is the fertility function.

We assume following equations held

$$\left[\frac{dQ_6}{dt}\right]_6 = x_6\left[\frac{dQ_2}{dt}\right]_6 - x_2\left[\frac{dQ_6}{dt}\right]_6$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_6 - x_4\left[\frac{dQ_2}{dt}\right]_6$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_6 - x_6\left[\frac{dQ_3}{dt}\right]_6$$

Where $x_6, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_6$ and $g_6(P_6)Q_6$

Now when you add the term $(b_6Q_6)$ with $g_6(P_6)(Q_6)$ we are infect see is that $(b_6Q_6)$ is credited incrementally part of $g_6(P_6)(Q_6)$ or instant nearly, when subtracted from $\frac{dQ_6}{dt}$ , it means that $g_6(P_6)(Q_6)$ is slowly eaten by $\frac{dQ_6}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_6 = \mu_6(a_6) + (f_6)_6(P_6)_6 + (f_2)_6(Q)_6 \quad (2)$$

$$n_6 = \beta_6(a_6)h_6(Q_6), \quad (3)$$

Where $\mu_6$the age- dependent debit rate whit is $\mu_6(a_6)$ 0 $\forall a > 0$; $(f_6)_6$ is the density- dependent death rate with $f_6(0) = 0$, $(f_6^6)_6(P_6)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_6)_6(P_6) = (k_6)_6$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_6(0) = 0, (f_2^6)_6(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_6(Q_6) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 6, h^6(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_6 + (\rho_a)_6 = -\mu_6(a_6) + (f_6)_6P_6) + \big((f_2)_6(Q_6)\big)(\rho_6)_6 \quad (4)$$

$$\rho_6(t_6, 0) = h_6(Q_6)\int_0^\infty \beta_6(a_6)\rho_6(t, a_6)da, \quad (5)$$

$$\frac{dQ_6}{dt} = -b_6Q_6 + g_6(P_6)Q_6 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_6$ and $Q_6$ of (4-6) satisfy

$$0 \ \rho_6(t, a_6) \ \rho_6^*(t, a_6),$$

$$0 \leq Q_6(t) \leq Q_6^*(t),$$

Where $\rho_6^*(t, a_6)$ and $Q_6^*(t)$, are solutions of

$$(\rho_t^*)_6 + (\rho_a^*)_6 = -(\mu_6(a_6) + (f_6)_6(P_6^*) + (f_2)_6 Q_6^*))\,\rho_6^*, \quad (7)$$

$$\rho_6^*(t, 0) = \int_0^\infty \beta_6(a_6)\,\rho_6^*(t, a_6)\,da_6 \quad (8)$$

$$\frac{dQ_6^*}{dt} = -b_6\,Q_6^* + g_6(P_6^*)Q_6^*, \quad (9)$$

With $P_6^*(t) = \int_0^\infty \rho_6^*(t, a_6)da_6$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_6^*(t)}{dt} = P_6^*(t)\big(r_6 - (f_6)(P_6^*) - (f_2)_6(Q^*)\big), \quad (10)$$

$$\frac{dQ_6^*(t)}{dt} = Q_6^*(t)(-b_6 + g_6(P_6^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_6(a)e^{-\int_0^a \mu_6(\tau)dt}\,e^{-ra}\,da_6 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_6^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(t)dt}\, da_6 > 1, \qquad (H1)$$

Which guarantees $r_6 > 0$.

*Remark.* If $\beta_6(a) = [\beta_0]_6 a_6 e^{-\alpha_6 a_6}, \mu_6(a_6) \equiv \mu_6$ then the solution of (12) is

$$r_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_6]_6(P_6) = 0, [f_2](Q_6) = \lambda_6 Q_6\ h_6(Q_6) = 1, \qquad and\ \ g_6(P_6) = c_6 P_6$$

$$[\rho_6]_6 + [\rho_a]_6 = -(\mu_6 + \lambda_6 Q_6)\rho_6$$

$$\rho_6(t,0) = \int_0^\infty [\beta_0]_6 a_6 e^{-\alpha_6 a_6} \rho_6(t, a_6)\, da_6$$

$$\frac{dQ_6}{dt} = -b_6 Q_6 + c_6 P_6 Q_6$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_6}{dt} = \omega_6 P_6 - \lambda_6 P_6 Q_6$$

$$\frac{dQ_6}{dt} = b_6 Q_6 - c_6 P_6 Q_6$$

Where $\omega_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$

## 58. THE EQUILIBRIUM $(\rho_0)_6(a_6), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_6(a_6), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_6 + (\rho_a)_6 = -\big(\mu_6(a_6) + (f_6)_6(P_6)\big)\rho,$$

$$\rho_6(t,0) = \int_0^\infty \beta_6(a_6)\rho_6(t,a_6)da_6 \qquad (13)$$

The solution of (63) has the form of

$$\rho_6(t,a_6) = Ce_6^{r(t-a_6)} - \int_0^a (\tau)d\tau - \int_0^6 (f_6)_6\big(P_6(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_6(t) = C_6 e^{r_6 t - \int_0^t (f_6)_6(P_6(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_6^{-6}(r)$ if $k_6 > r$. This is equivalent to

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{(k_6)_6 a_6} da_6 < 1.$$

In fact, if we define

$$R_6(s_6) = \int_0^\infty \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau)d\tau} e^{-s_6 a_6}\, da_6$$

$R(s)$ is a decreasing function. $R_6(r_6) = 1$ And $R_6(k_6)_6 < 1$ imply $r_6 < (k_6)_6$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_6(a_6), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{k_6 a_6} da_6$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_6 \equiv \int_0^\infty (\rho_0)_{66}(a_6)da_6$ *satisfies*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{-(f_6)_6(P_0)a_6} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_6(a_6), 0)$

Let $u_6(t, a_6) = \rho(t, a_6) - \rho_0(a_6)$ Then

$$N_6(t) \equiv \int_0^\infty u_6(t, a_6)da_6 = \int_0^\infty \rho_6(t, a_6)da_6 - \int_0^\infty (\rho_0)_6(a_6)da = P_6(t) - P_0)_6$$

And

$$(u_6)_6 + (u_a)_6 = (\rho_t)_6 + (\rho_a)_6 - (\rho_0^6)_6(a_6)$$

$$= -\mu_6(a_6) + (f_6)_6(P_6) + (f_2)_6(Q_6))\,(u_6 + (\rho_0)_6) + \mu_6(a_6)(\rho_0)_6 + (f_6)_6(P_0)_6(\rho_0)_6$$

$$= -\mu_6(a_6)u_6 - (f_6)_6(P_6) + (f_2)_6(Q_6))\,u_6 - (\rho_0)_6\,(f_6)_6(P_6) - (f_6)_6(P_0)_6$$

$$+ (f_2)_6(Q_6)), \qquad (15)$$

Where $(\rho_0^6)_6(a_6)$ denotes $d(\rho_0)_6(a_6)/da_6$

Expansion $(f_6)_6(P_6)$ and $(f_2)_6(Q_6)$ at $(P_0)_6$ and 0 yield respectively

$$(f_6)_6(P_6) = (f_6)_6((P_0)_6 + N_6)$$

$$= (f_6)_6(P_0)_6 + (f_6^6)_6(P_0)_6 N_6 + o_6(N_6), \qquad (16)$$

$$(f_2)_6(Q_6)$$

$$= (f_2^6)_6(0)Q + o_6(Q_6) \quad (17)$$

Write $f_6(P_0) = f_{60}, f_6^{'}(P_0) = f_{60}^{'}$, and $f_2^{'}(0) = f_{20}^{'}$ and substitute (16) and (17) into (15) to obtain

$$(u_6)_6 + (u_a)_6 \approx -(\mu_6(a_6) + (f_{60})_6)u - (f_{60}^6)_6(\rho_0)_6(N_6) - (f_{20}^6)_6(\rho_0)_6(Q_6)$$

$$u_6(t,0) = \rho_6(t,0)(\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a) + (\rho_0)_6(a_6)da_6 - (\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a_6)da_6 + (h_6(Q_6) - 1)(\rho_0)_6(0) \quad (18)$$

Expanding

$$h_6(Q_6) = 1 + (h^6)_6(0)Q_6 + 0(Q_6) \equiv 1 + h_0^{'}Q_6 + 0Q_6$$

And substituting it into (18), we arrive at

$$u_6(t,0) \approx \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + (h_{a_6}^{'})(\rho_0)_6Q_6$$

The linearization is the following :

$$[u_6] + [u_6] = -(\mu_6(a_6) + [t_{60}]_6u_6 - [\alpha_6]_6(a_6)N_6 - [\alpha_2]_6(a_6)Q_6$$

$$u_6(t,0) = \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + [\xi_6]Q_6 \quad (19)$$

$$\frac{dQ_6}{dt} = -b_6Q_6$$

where

$$[\alpha_6]_6(a_6) = [t_6^6 a]_6[p, o]_6(a_6)d2(a_6) = [f_{20}^6]_6(\rho\, o\,)_6(a_6)\xi = [\mathrm{h\ o}^6]_6[\rho \mathrm{o}]_6(\mathrm{o}) \text{ and } \mathrm{b}_6 = \mathrm{b}_6 - g_6(P_6)$$

If $g_6(P_0) > b$, clearly, $(\rho_0)_6(a_6), 0)$ is unstable and the predator population is persistent. Assume $g_6(P_0)_6 < b$.

Then, $Q_6(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_6 + u_a = -(\mu_6(a_6) + (f_{60})_6)u_6 - (\alpha_6)_6(a_6)N_6$$

$$= -(\mu_6(a_6) + f_{60})u_6 - \alpha_6(a_6)\int_0^\infty u_6(t, a_6)da_6$$

$$u_6(t, 0) = \int_0^\infty \beta_6\,(a_6)u_6(t, a_6)da_6$$

Using the method developed by Cushing [2, 3], we suppose

$$u_6(t, a_6) = \tilde{\mathrm{u}}_6(a_6)^{e^{\gamma(t-a_6)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_6}\, da_6$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_6)$ gives,

$$\tilde{\mathrm{u}}_6(a_6) = e^{-\int_0^a \mu_6(\tau_6)d\tau - (f_{60})_6 a}\left(\tilde{\mathrm{u}}_6(0) - w\int_0^a e^{\int_0^\tau \mu_6(s_6)ds + (f_{60})}\tau_6\,(\alpha_6)_6(\tau_6)e^{\gamma_6\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_6(a_6)e^{-\gamma_6 a_6 - (f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6)d\tau}\right)\tilde{\mathrm{u}}_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a(\alpha_6)_6(\tau_6)e^{-\gamma_6\tau_6}e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds}d\tau_6\,da_6$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}da\right)(\tilde{u}_6(0))_6$$

$$+\left(-6-\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds+y\tau}\alpha_6(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(t_6)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_6\,(a_6)Y_6(\gamma_6,a_6)da_6\right)u_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6=0,$$

(24)

$$\left(\int_0^{\infty}Y_6(\gamma_6,a_6)da_6\right)\tilde{u}_6(0)$$

$$+\left(-1-\int_0^{\infty}Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6\right)w=0.$$

Equqtion (20) has a solution of the form of $u_6(t,a_6)=u_6(a_6)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_6(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) Y^{-6}(\gamma_6, a_6) d\tau_6$. We obtain the characteristic equation

$$f_{60}^{l}{}_6 \left( \int_0^\infty Y_6(\gamma_6, a_6) da_6 \right) \left( \int_0^\infty \beta_6(a_6) Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$+ \left( 1 - \int_0^\infty \beta_6(a_6) Y_6(\gamma_6, a_6) da \right) \left( 1 + (f_{60}^{l})_6 \int_0^\infty Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$= 0. \qquad (25)$$

Since

$$(\rho_0)_6(a_6) = (\rho_0)_6(0) e^{-\int_0^a \mu_6(\tau_6) d\tau_6 - (f_{60})_6 a_6}$$

$$I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) e^{\gamma_6 \tau_6 + (f_{60})_6 \tau + \int_0^\tau \mu_6(s_6) ds_6} d\tau$$

$$= \frac{(\rho_0)_6(0)}{\gamma_6} (e^{\gamma_6 a_6} - 1)$$

Let

$$E_6(a_6) = e^{-(f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6) d\tau_6}$$

This relationship is cardinal one. It is useful in quantum computation and other black hole dynamics problems

**Quantum space and quantum time**

(3237) **Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.**

(3238) **In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.**

(3239) **They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.**

(3240) **When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.**

(3241) **Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.**

**The Planck scale**

(3242) **Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".**

(3243) **Here ħ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.**

(3244) **The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.**

(3245) **Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)**

(3246) **Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.**

## Physical applications

**Some phenomena where quantum gravity plays a major role are the following.**

**The microscopic structure of spacetime**

(3247) **If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.**

(3248) **Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length**

(3249) **Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.**

(3250) **A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.**

**Early cosmology**

(3251) **According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.**

(3252) **However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.**

(3253) **A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.**

(3254) **Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".**

(3255) **One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.**

**Black holes**

(3256) **Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.**

(3257) **First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.**

(3258) **Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.**

(3259) **What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.**

(3260) **Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.**

**Astrophysical effects**

(3261) **Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.**

(3262) **An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.**

(3263) **The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.**

## Main approaches

(3264) **The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.**

(3265) **Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.**

(3266) **The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.**

## String theory

(3267) **String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.**

(3268) **Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.**

(3269) **String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.**

(3270) **Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.**

(3271) **However, when summed the entire series appears to be (=) divergent.**

(3272) **In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.**

(3273) **Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.**

(3274) **Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.**

(3275) **In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.**

## Loop quantum gravity and spin foams

(3276) **Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.**

(3277) **The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity.**

(3278) **Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime.**

(3279) **The granular properties of space can be (=) explicitly computed.**

(3280) **In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory.**

(3281) **The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks.**

(3282) **Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure.**

(3283) **The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit.**

(3284) **Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism.**

(3285) **This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism.**

## Noncommutative geometry

(3286) **Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime.**

(3287) **Quantum mechanics teaches us that dynamical quantities are (=) noncommutative -- meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed.**

(3288) **It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry.**

(3289) **A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces.**

(3290) **Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces.**

(3291) **This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity.**

**Others**

(3292) **Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays**

(3293) **. Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection.**

(3294) **Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible.**

(3295) **The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.**

**Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016**

(3296) **These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field.**

(3297) **The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets.**

(3298) **Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state.**

(3299) **Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands.**

(3300) **The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI**

(3301) **The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9].**

(3302) **This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive.**

(3303) **More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.**

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

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| (3304) | **Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.** |
| (3305) | **Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.** |
| (3306) | **The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.** |
| (3307) | **This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.** |
| (3308) | **The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.** |
| (3309) | **This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.** |
| (3310) | **For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.** |
| (3311) | **The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society** |

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

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| (3312) | **Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.** |
| (3313) | **We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.** |
| (3314) | **When this is true—as is often the case—the system serves as (=) a "bath" for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called thermalization.** |
| (3315) | **However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.** |
| (3316) | **Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1).** |
| (3317) | **Bringing them to thermal equilibrium requires (e) an external bath.** |

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| (3318) **But so far, there has been no rigorous proof of (e) the existence of such many-body localization. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption.** |
| (3319) **The concept of localization was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations).** |
| (3320) **The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century.** |
| (3321) **But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information.** |
| (3322) **Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment.** |
| (3323) **For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood.** |
| (3324) **It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7].** |
| (3325) **Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature.** |
| (3326) **One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.** |
| (3327) **However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].** |
| (3328) **Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.** |
| (3329) **Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.** |
| (3330) **As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked.** |
| (3331) **His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice.** |
| (3332) **An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site.** |
| (3333) **This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself).** |

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| (3334) **This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition.** |
| (3335) **However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime:** |
| (3336) **These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects.** |
| (3337) **It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies).** |
| (3338) **This is a realistic scenario for (e) all known systems.** |
| (3339) **On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects.** |
| (3340) **On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other.** |
| (3341) **No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true.** |
| (3342) **His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found.** |
| (3343) **Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization?** |
| (3344) **Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects?** |
| (3345) **Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions.** |
| (3346) **So in those systems, the quantum phase transition does indeed exist.** |
| (3347) **But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension.** |
| (3348) **Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters.** |

**Focus: Biological Cells Form Electric Circuits July 1, 2016• Physics 9, 75 Cells that are electrically active and that also produce light for easy voltage monitoring could lead to new studies of heart arrhythmias and possibly bio-computing. Figure caption expand figure H. M. McNamara et al., Phys. Rev. X (2016) Cell circuit. Waves of electrical activity propagate in a loop within a collection of cells engineered with electrical and optical properties (see video below).Some sentences, pictures and videos are deleted due to constraints on space.**

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| (3349) **The human heartbeat is produced by (e) electrical pulses that propagate through (e&eb) cardiac tissue, causing (eb) rhythmic muscle contraction.** |
| (3350) **Researchers have previously engineered cells to create (eb) an artificial tissue capable of producing (eb) coordinated electrical activity, and now a team has added the ability to monitor (e&eb) their electrical state by (e) detecting fluorescent emission.** |
| (3351) **They have also fashioned the cells into (e&eb) "living circuits" that might act as (=) model systems for studying heart behavior.** |
| (3352) **Constructing electrically active tissues from their basic building blocks by (e) genetically engineering cells allows (eb) researchers to study a "clean" system similar to (e) the more complicated cardiac tissue.** |
| (3353) **These systems let them investigate, for example, the electrical signals that lead to (eb) the potentially fatal irregular heartbeat called arrhythmia.** |
| (3354) **Previously researchers have introduced (eb) ion channels—proteins that conduct ions through (e&eb) cell membranes—into human kidney cells to make (eb) electrically "excitable" tissues [1].** |
| (3355) **Adam E. Cohen of Harvard University and his co-workers decided to add more ingredients to such cells, so that the electrical activity could be (=) light-activated and easily visualized.** |
| (3356) **They incorporated four additional gene products into (e&eb) human embryonic kidney (HEK) cells.** |
| (3357) **Two were ordinary ion channels; the other two were (=) a blue-light-activated ion channel (a light-sensitive switch) and (e&eb) a protein that fluoresces with (e&eb) an intensity that depends on (=) the voltage across the membrane in which it is embedded.** |
| (3358) **Light activation of electrically excitable cells has been shown before [2]. But Cohen and colleagues have now included a convenient "readout" mechanism (the fluorescent protein) and have used patterning methods to shape (e&eb) the cell cultures into "tissue circuits."** |
| (3359) **They grew rings of cells a few millimeters in diameter that could act as (e&eb) oscillators that support (eb) circulating waves of electrical excitation.** |
| (3360) **These waves arise from (e) the collective behavior of the cells: a pulse produced by (e) one cell can trigger another from (e) its neighbors, eventually producing (eb) a coherent wave of excitation traveling through (e&eb) the tissue.** |
| (3361) **It's this coordination, in a repeating cycle of excitation and quiescence, which gives rise to (eb) rhythmic activity in cardiac tissue.** |

**H. M. McNamara et al., Phys. Rev. X (2016)**

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| (3362) **A pattern of blue light (not visible) illuminates these cells, defining (eb) a rectangular loop around which electrical waves can propagate.** |
| (3363) **At about 0:07 the rectangle begins to shrink in (eb) the vertical direction, causing t(eb) he cycle time to gradually decrease. (In other experiments, the cells were grown in (eb) fixed, ring-shaped patterns.)** |
| (3364) **The ring of engineered HEK cells acted like an electronic "ring oscillator," a device with a characteristic frequency to (e&eb) its circulating electrical signal.** |
| (3365) **The electrical waves in a ring oscillator can propagate (e&eb) in either direction around the ring, enabling (eb) it to serve as a binary memory unit.** |

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| (3366) **Cohen and colleagues found that in their HEK tissues, these circulating waves remained (eb) stable for more than two hours, with (e&eb) the pulses completing a few circuits every second.** |
| (3367) **The electrical signals were easily visualized as (=) waves of fluorescent emission.** |
| (3368) **The researchers say that these tissues could be used to (e) understand and predict the behavior of real cardiac cells, especially arrhythmias.** |
| (3369) **"In real heart, if you pace the tissue fast enough, *(eb) it starts skipping beats and eventually becomes (=) arrhythmic," says Cohen.** |
| (3370) **This state is a well-known precursor to (e) sudden cardiac death and is associated with (e&eb) the breakup of the usual regular waves of electrical excitation.** |
| (3371) **Some drugs also have the dangerous side effect of inducing (eb) arrhythmia, and Cohen says that the new system might offer a quick and convenient way of screening for such hazards. The work "is really very elegant and deserves attention," says physiologist Gil Bub of Oxford University in the UK. "It opens new research questions for people like me who are interested in wave dynamics in excitable systems."** |
| (3372) **There might also be practical uses for information processing in living tissues. "By combining patterns of cell growth and (e&eb) patterns of ion channel expression, we think we can build diodes, amplifiers, and other simple logic elements," says Cohen.** |
| (3373) **Andrew Adamatzky of the University of the West of England in Bristol, a specialist in "unconventional computing," foresees circuits that perform computations through (e&eb) the interactions of electrical waves. "This research might lead to fabrication of implantable bio-computers," he says. This research is published in Physical Review X. –Philip Ball Philip Ball is a freelance science writer in London and author of Invisible: The Dangerous Allure of the Unseen (University of Chicago Press, 2014).** |

**Optically Controlled Oscillators in an Engineered Bioelectric Tissue Harold M. McNamara, Hongkang Zhang, Christopher A. Werley, and Adam E. Cohen Phys. Rev. X 6, 031001 – Published 1 July 2016**

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| (3374) **Complex electrical dynamics in excitable tissues occur throughout (e&eb) biology, but the roles of individual ion channels can be difficult to determine due to (e) the complex nonlinear interactions in (eb) native tissue.** |
| (3375) **Here, authors ask whether they can engineer a tissue capable of (e) basic information storage and processing, where (e) all functional components are known and well understood.** |
| (3376) **Authors develop a cell line with (e&eb) four transgenic components: two to enable (eb) collective propagation of electrical waves and (e&eb) two to enable optical perturbation and (e&eb) optical readout of membrane potential.** |
| (3377) **They pattern the cell growth to define (eb) simple cellular ring oscillators that run stably for >2 h (∼104 cycles) and that can store data encoded in (eb) the direction of electrical circulation.** |
| (3378) **Using (e) patterned optogenetic stimulation authors probe the biophysical attributes of (e) this synthetic excitable tissue in detail, including (e) dispersion relations, (e&eb) curvature-dependent wave front propagation, (e&eb) electrotonic coupling, and (e&eb) boundary effects.** |
| (3379) **Authors then apply the biophysical characterization to (e) develop an optically reconfigurable bioelectric oscillator.** |
| (3380) **These results demonstrate the feasibility of engineering bioelectric tissues capable of (e&eb) complex information processing with optical input and output. Received 12 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevX.6.031001 This article is available under the terms of the** |

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**Scientists detect radio emission from a nearby brown dwarf July 6, 2016 by Tomasz Nowakowski report**

(3381) **K2 and Spitzer images of W0607+24: The position of W0607+24 is marked by (e) a red circle. SDSS J060738.06+242939.5, marked by (e) an orange circle, are the brightest star in the K2 image.**

(3382) **The three faint (r = 18.3 − 19.6) stars marked by cyan circles, contribute significant flux to (e) the K2 measurement but are (-) negligible at Spitzer's longer wavelength and (e &êb) higher resolution. Other stars with r < 20 in SDSS are marked with blue circles. Credit: Gizis et al., 2016.**

(3383) **(Phys.org)—Located some 23 light years away, the brown dwarf designated WISEP J060738.65+242953.4 or W0607+24 for short, turns out to be (=) a source of radio emission.**

(3384) **According to a research paper published July 4, this substellar object showcases (eb) quiescent radio emission, making it one of (e) the most radio-faint, ultra-cool dwarfs yet detected.**

(3385) **Brown dwarfs like W0607+24 are objects that are too large to be called planets and too small to be stars. With a mass below that necessary to maintain hydrogen-burning nuclear fusion reactions, they are much cooler and dimmer than (eb) main sequence stars.**

(3386) **These objects could be also sources of (e) radio emission, but the cause of this process is still not completely understood.**

(3387) **Due to its proximity, W0607+24 is (=) an excellent target for studying the nature of (e) brown dwarfs.**

(3388) **It is the nearest known late-L dwarf in the northern hemisphere, and the third-nearest in (=) the whole sky, (eb) allowing researchers to study phenomena occurring in brown dwarfs.**

(3389) **That's why a team of astronomers, led by John Gizis of the University of (e)Delaware, Newark, carried out a series of mid-infrared, radio, and optical observations of W0607+24 to study any activity in this curious substellar object.**

(3390) **Motivated by the results from NASA's prolonged Kepler mission, named K2, which obtained (eb) long time-series photometry of W0607+24; they observed this object simultaneously with the Karl G. Jansky Very Large Array (VLA) in New Mexico and NASA's Spitzer Space Telescope.**

(3391) **The most important results were provided by VLA, as the team found that W0607+24 is (=) a source of quiescent, non-bursting, radio emission.**

(3392) **The scientists noted that the discovery of this process in this nearby object could be (=) crucial for our understanding of this process in (eb) ultracool dwarfs in general.**

(3393) **"W0607+24 is a valuable test case for (e) understanding the mechanisms of radio emission in ultracool dwarfs.**

(3394) **Some of the radio-detected objects emit (eb) rapid, intense bursts with high circular polarization, often found to occur periodically and attributed to coherent emission due to (e) the electron cyclotron maser instability," the paper reads.**

(3395) **The astronomers also suggested that the relative radio faintness of W0607+24 may be due to (e) its orientation if (e) the emission is concentrated at the poles.**

(3396) **However, follow-up observations are needed to better characterize (eb) the detected radio emission and to eliminate (e) all the uncertainties.**

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| (3397) **The team also hopes for more significant findings when (e) studying this nearby substellar object in the future.** |
| (3398) **"As the nearest known northern hemisphere late-L dwarf, this offers many possibilities for future studies. For example, it could serve as (=) an useful standard for (e) line-broadening and (e&eb) polarization studies.** |
| (3399) **If there is a planetary system aligned with (e&eb) (the primary's rotation axis, then we would expect very little radial velocity signal but a strong astrometric signal," the researchers concluded.** |
| (3400) **Moreover, the team plans to monitor (e&eb) additional late-M and L dwarfs in other K2 campaigns, expecting even more promising results.** |
| (3401) **They are more distant than W0607+24, but because they are warmer, (eb) some are brighter and higher signal-to-noise than that described in this paper.** |
| (3402) **In addition to detecting the radio emission, (eb)the scientists also estimated the age, mass and radius of W0607+24** |
| (3403) **. According to the study, the object is less than two billion years old, has (e) a maximum mass of about 0.055 solar masses and a radius approximately of 0.1 solar radii.** |
| (3404) **More information: WISEP J060738.65+242953.4: A Nearby. Pole-On L8 Brown Dwarf with Radio Emission, arXiv: 1607.00943 [astro-ph.SR] arxiv.org/abs/1607.00943** |
| (3405) **Authors present a simultaneous, multi-wavelength campaign targeting the nearby (7.2 pc) L8/L9 (optical/near-infrared) dwarf WISEP J060738.65+242953.4 in (eb) the mid-infrared, radio, and optical.** |
| (3406) **Spitzer Space Telescope observations show no (e) variability at the 0.2% level over 10 hours each in the 3.6 and 4.5 micron bands.** |
| (3407) **Kepler K2 monitoring over 36 days in Campaign 0 rules out (e) stable periodic signals in (eb) the optical with amplitudes great than 1.5% and periods between 1.5 hours and 2 days.** |
| (3408) **Non-simultaneous Gemini optical spectroscopy detects lithium, constraining (e) this L dwarf to be less than ~2 Gyr old, but no Balmer emission is observed.** |
| (3409) **The low measured projected rotation velocity (v sin i < 6 km/s) and lack of variability are (=) very unusual compared to other brown dwarfs, and authors argue that this substellar object is (=) likely viewed pole-on.** |
| (3410) **Authors detect quiescent (non-bursting) radio emission with (e&eb) the VLA.** |
| (3411) **Amongst radio detected L and T dwarfs, it has (e) the lowest observed L_nu and the lowest v sin i.** |
| (3412) **They discuss the implications of pole-on detection for (e) various proposed radio emission scenarios. © 2016 Phys.org** |

**A new look at the galaxy-shaping power of black holes July 6, 2016**

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| (3413) **It wasn't until the advent of X-ray astronomy that the full picture of (e) the galaxy-shaping power of black holes began to emerge (eb) with the ability to see plasma.** |
| (3414) **Data from a now-defunct X-ray satellite is providing (eb) new insights into the complex tug-of-war between galaxies, the hot plasma that surrounds them, and the giant black holes that lurk in their centres.** |

(3415) **Launched from Japan on February 17, 2016, the Japanese space agency (JAXA) Hitomi X-ray Observatory functioned for just over a month before contact was lost and (e&eb) the craft disintegrated.**

(3416) **But the data obtained during those few weeks was enough to paint (eb) a startling new picture of the dynamic forces at work within galaxies.**

(3417) **New research, published in the journal Nature today, reveals data that shows just how important the giant black holes in galactic centres are to the evolution (eb) of the galaxies as a whole.**

(3418) **"We think that supermassive black holes act(e&eb) like thermostats," said Brian McNamara, University Research Chair in Astrophysics at the University of Waterloo. "They regulate (e&eb) the growth of galaxies."**

**Champagne bubbles of plasma**

(3419) **During its brief life, the Hitomi satellite collected X-ray data from (e) the core of the Perseus cluster, an enormous gravitationally-bound grouping of (e) hundreds of galaxies.**

(3420) **Located some 240 million light years from earth, the Perseus cluster is one of (e) the largest known structures in the universe.**

(3421) **The cluster includes not only the ordinary matter that makes up the galaxies, but an (e) "atmosphere" of hot plasma with a temperature of tens of millions of degrees, as well as a halo of (e) invisible dark matter.**

(3422) **Earlier studies, going back to the 1960s, have shown that each of the galaxies in the cluster - and indeed most galaxies - likely contains (e) a supermassive black hole in its centre, an object 100 million to more than ten billion times as massive as our sun**

**New look at the galaxy-shaping power of black holes**

(3423) **Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving.**

(3424) **"These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science.**

(3425) **"Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy."**

(3426) **This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return.**

(3427) **The remaining matter gets absorbed into the black hole, adding to (e) its mass.**

(3428) **The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne.**

(3429) **The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution.**

(3430) **Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The**

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| **answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form.** |
| (3431) **"Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara.** |
| (3432) **"It creates (eb) these bubbles, and the bubbles keep the plasma hot.** |
| (3433) **That's what prevents galaxies from (e) becoming even bigger than they are now."** |
| (3434) **Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.** |
| (3435) **In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.** |
| (3436) **Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.** |
| (3437) **"It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.** |
| (3438) **"It's like nature's thermostat, which keeps these galaxies from (e) growing.** |
| (3439) **If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."** |

**A new look at the galaxy-shaping power of black holes**

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| (3440) **An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the** |
| (3441) **McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.** |
| (3442) **Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.** |

**Doomed satellite missions**

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| (3443) **Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.** |
| (3444) **At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.** |
| (3445) **Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a** |

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| **leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.** |
| (3446) **The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.** |
| (3447) **These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.** |
| (3448) **"But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people** |
| (3449) **. It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."** |
| (3450) **The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.** |
| (3451) **"This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website** |

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

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| (3452) **Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.** |
| (3453) **They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.** |
| (3454) **However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.** |
| (3455) **Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.** |
| (3456) **X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well.** |
| (3457) **The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma.** |
| (3458) **These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6.** |
| (3459) **Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus.** |

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| (3460) | **A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate.** |
| (3461) | **Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure.** |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (3462) | **This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy.** |
| (3463) | **Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole.** |
| (3464) | **But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy.** |
| (3465) | **A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole.** |
| (3466) | **It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth.** |
| (3467) | **"There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA).** |
| (3468) | **This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies.** |
| (3469) | **However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated.** |
| (3470) | **In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1.** |
| (3471) | **We know dark matter exists only from (e) its gravitational effects.** |
| (3472) | **It holds together galaxies and (e&eb) galaxy clusters.** |
| (3473) | **Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years.** |
| (3474) | **To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies.** |
| (3475) | **They used star motions as tracers to (e) weigh the galaxies' central black holes.** |
| (3476) | **X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto.** |
| (3477) | **They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone.** |
| (3478) | **This connection is likely to be related to (e&eb) how elliptical galaxies grow.** |

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| (3479) **An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together.** |
| (3480) **Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole.** |
| (3481) **"In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp** |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (3482) **The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the …** |
| (3483) **Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness?** |
| (3484) **Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness.** |
| (3485) **"Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain."** |
| (3486) **Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness.** |
| (3487) **"With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication."** |
| (3488) **Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain** |

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| **changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories.** |
| (3489) **The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen.** |
| (3490) **In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer.** |
| (3491) **Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants.** |
| (3492) **These were treated as (=) "aware" and "unaware" trials, respectively.** |
| (3493) **Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it.** |
| (3494) **The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another.** |
| (3495) **Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion.** |
| (3496) **No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness.** |
| (3497) **"We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said.** |
| (3498) **"Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness."** |
| (3499) **The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner.** |
| (3500) **Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain.** |
| (3501) **"We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said.** |
| (3502) **"This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients** |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

(3503) **Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter.**

(3504) **As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined.**

(3505) **Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions.**

(3506) **"Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.**

(3507) **The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.**

(3508) **However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.**

(3509) **The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.**

(3510) **Noda and his colleagues instead studied the change in emissions over (e&eb) time.**

(3511) **By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.**

(3512) **Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).**

(3513) **The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.**

(3514) **Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.**

(3515) **"The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp**

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(3516) **The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking "palm tree" organization.**

(3517) **The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society**

(3518) **(Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible?**

(3519) **With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem.**

(3520) **This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other.**

(3521) **According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible.**

(3522) **This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges.**

(3523) **In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century.**

(3524) **Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching.**

(3525) **In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom.**

(3526) **Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth.**

**Spheres on trees**

(3527) **Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration.**

(3528) **They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement.**

(3529) **Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum.**

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| (3530) **By visualizing the problem in this way, the researchers noticed that these particular graphs don't (e) have lots of branches extending from lots of (e) other branches.** |
| (3531) **Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure.** |
| (3532) **The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network.** |
| (3533) **As a result, it provides (eb) an important clue for finding the global minimum.** |
| (3534) **It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum.** |
| (3535) **This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large.** |

**Thompson's problem 2**

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| (3536) **By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society** |
| (3537) **The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima.** |
| (3538) **The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node.** |

**Implications for Thomson's and Smale's problems**

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| (3539) **Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before.** |
| (3540) **"Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around."** |
| (3541) **As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem.** |
| (3542) **"I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view."** |
| (3543) **In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges.** |
| (3544) **From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear.** |

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| (3545) | **"While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK.** |
| (3546) | **"For larger numbers, there will likely be (=) multiple funnels.** |
| (3547) | **Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum."** |
| (3548) | **Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.** |
| (3549) | **"The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.** |
| (3550) | **"Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org** |

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**Introduction:**

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price, University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss.**

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| **(577)** | **Black Holes are regions of (e) space in which gravitational fields are so strong that no particle or signal can escape (eb) the pull of gravity.** |
| **(578)** | **The boundary of this no-escape region is (=) called the event horizon, since distant observers outside the black hole cannot (e) see (cannot get light from) events inside.** |
| **(579)** | **Although the fundamental possibility of such an object exists within (eb) Newton's classical theory of gravitation, Einstein's theory of gravity makes (eb)black holes inevitable under some circumstances.** |
| **(580)** | **Prior to the early 1960s, black holes seemed to be (=) only an interesting theoretical concept with no (e) astrophysical plausibility, but with (e&eb) the discovery of quasars in 1963 it became clear that very exotic astrophysical objects could exist.** |
| **(581)** | **Nowadays it is taken for granted that black holes do exist in (eb) at least two different forms.** |
| **(582)** | **Stellar mass black holes are (=) the endpoint of the death of some stars, and supermassive black holes are the result of (e) coalescences in the centers of most galaxies, including our own.** |
| **(583)** | **No signal can propagate from (e) inside a black hole, but the gravitational influence of a black hole is (=) always present.** |

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| (584) | This influence does not (e) propagate out of the hole; it is permanently present outside, and depends only on (=) the total amount of mass, angular momentum, and electric charge that have gone into forming (eb) the hole. |
| (585) | Black holes can be detected through the influence of this (e&eb) strong gravity on the surroundings just outside the hole. |
| (586) | In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them. |
| (587) | Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors. |

**Classical vs. relativistic black holes**

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| (588) | Something like a black hole exists within (eb) Newton's classical theory of gravity. |
| (589) | In that theory, an energy argument tells us that there is (=) an escape velocity $v_{esc}=\sqrt{2GM/R}$ from (e) the surface of any spherical object of (e) mass M and radius R |
| (590) | If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity. |
| (591) | Thus the condition for such an "unseeable" object is $R<2GM/c^2$ |
| (592) | In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape. |
| (593) | This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity. |
| (594) | To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it. |
| (595) | The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c . |
| (596) | This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed. |
| (597) | An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity. |
| (598) | Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable. |

**Stationary black holes**

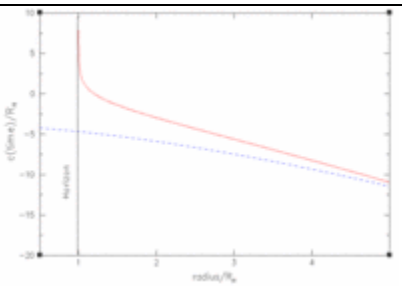
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| (599) | In the Newtonian theory, gravity is described by (e) the potential $\Phi$. |
| (600) | Inside a spherical object the form of (e) $\Phi(r)$ depends on (=) the interior structure, but in the vacuum outside (e) matter the potential $-GM/r$ depends only on (=) the interior mass. |
| (601) | Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object. |

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| **(602) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward.** |
| **(603) This horizon radius RH=2GM/c2 is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory.** |
| **(604) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution.** |
| **(605) Radius RHhere actually means that the area of the event horizon is (=) 4πR2H.** |
| **(606) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution.** |
| **(607) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity.** |
| **(608) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own.** |
| **(609) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole.** |
| **(610) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density.** |
| **(611) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed.** |
| **(612) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole.** |
| **(613) In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter.** |
| **(614) Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum.** |
| **(615) In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior.** |
| **(616) The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object.** |
| **(617) Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum.** |
| **(618) These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.]** |
| **(619) All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time.** |
| **(620) But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?"** |
| **(621) The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point.** |

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| **(622) One of these coordinates is (=) called the "coordinate time."** |
| **(623) Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate.** |



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| **(624) Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole** |
| **(625) At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole.** |
| **(626) Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations.** |
| **(627) For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree** |
| **(628) An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events.** |
| **(629) The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving.** |
| **(630) Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle.** |
| **(631) The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return.** |
| **(632) But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon.** |
| **(633) Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future.** |
| **(634) The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?"** |

**Black hole parameters**

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| **(635) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin).** |
| **(636) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun.** |
| **(637) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits.** |
| **(638) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form.** |
| **(639) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds.** |
| **(640) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2.** |
| **(641) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed.** |
| **(642) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer.** |
| **(643) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2)** |
| **(644) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials.** |
| **(645) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable** |
| **(646) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite.** |
| **(647) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole.** |
| **(648) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars.** |
| **(649) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity.** |
| **(650) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity.** |
| **(651) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite.** |
| **(652) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity.** |
| **(653) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum.** |

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| **(654) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship.** |
| **(655) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= (4πG2/c4)[(M+M2−a2−−−−−−−√)2+a2].** |
| **(656) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored.** |

**Dynamical black holes**

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| **(657) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape.** |
| **(658) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, a horizon can be born.** |
| **(659) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid.** |
| **(660) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward.** |
| **(661) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations.** |
| **(662) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes.** |
| **(663) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3.** |
| **(664) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/** |
| **(665) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers.** |
| **(666) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission.** |
| **(667) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.** |
| **(668) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.** |
| **(669) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.** |
| **(670) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.** |

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| **(671) The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature 6×10−8 (Mo/M) K.** |
| **(672) Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.** |

Then $[Y_6]_6[\gamma_6]_6, a_6) = E_6(a_6)e^{-\gamma_6 a_6}$. by a simple calculation,

$$\int_0^\infty \beta_6(a_6)Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma}\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}\ da\right);$$

$$\int_0^\infty Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma_6}\left(\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}\ da_6\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_6(0)}{\gamma_6}[f_{60}^6]\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}da_6\ (\int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6)$$

$$+\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right).$$

$$\left(1 + \frac{(\rho_0)_6(0)}{\gamma_6}(f_{60}^{'})_6\left(\int_0^\infty E_6(a_6)da_6 - \int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}da_6\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right)(1 + \frac{(f_{60}^{'})_6}{\gamma_6}(\rho_0)_6(0)\int_0^\infty E_6(a_6)da_6) = 0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1 - \int_0^\infty \beta_6(a_6) Y_6(\gamma_6 a_6) da_6 = 1 - \int_0^\infty \beta_6(a_6) E_6(a_6) e^{-\gamma_6 a_6} da_6 \neq 0,$$

Since $\int_0^\infty \beta_6(a_6) Y_6(\gamma_6 a_6) \int_0^a \alpha_6(\tau_6) Y^{-6}(\gamma_6, \tau_6) d\tau_6 da_6 > 0$. Therefore, the characteristic solution is

$$\gamma_6 = -(f_{60}^{'})_6 (\rho_0)_6 \int_0^\infty E_6(a_6) da_6.$$

Since $f_6'(P_0)_6 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_6), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE OF DEBIT TYPES AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_6), 0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_6 + (\mu_6(a_6) + [f_6]_6(P_6^*) + [f_2]_6(Q_6^*))P_6^* = 0, \quad (26)$$

$$p_6^*(0) = h_6(Q_6^*) \int_0^\infty \beta_6(a_6) p_6^*(a_J) p_6^*(a_6) da_6 \quad (27)$$

$g_6(P_6^*) - b_6 = 0,$ (28)

Where $P_6^*$=$\int_0^\infty p_6^*(a_6)\, da_6$

Equation (28) yields $P^* = g^{-6}(b)$,

From (26),

$p_6^*(a) = p_6^*(0)e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6}$

Substituting into (27), we need to solve

$$\left(1 - h_6(Q_6^*)\int_0^\infty \beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6} da_6\right) p_6^*(0) = 0$$

For $Q_6^*$

Fix $p_6^*(0)$ and $P_6^*$ Let

$$H_6(z_6) = 1 - h_6(z_6)\int_0^\infty \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6} da_{,6}$$

And assume that

$$[k_3]_6 \int_0^\infty \beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(g_6^{-6}(b_6))a_6 - [k_2]_6 a_6}\, da_6 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_6(z_6)] > 0.$$

Clearly,

$$\frac{dH_6}{dz_6} = h_6(z)([f_2^{'}]_6(z_6))\int_0^{\infty} a_6\beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - f_6(P_6^*)a_6 - f_2(z_6)a_6}\, da_6$$

$-h^{'}(z_6)\int_0^{\infty}\beta_6(a_6)\, e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6}\quad da_6$

By the assumptions on $\left[f_2^6\right]_6 and\ [h^6]^6 dH_6/dz > 0$. Hence, if

$$H_6(0) = 1 - \int_0^{\infty}\beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6}da_6 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_6(0) < 0$ follows if $g_6^{-}(b_6) < [f_6^{-6}]_6(r_6)$, where $r_6$ is the solution of (12).

Indeed, since

$$R_6(s_6) = \int_0^{\infty}\beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 e^{-y}6 a}da_6$$

Is a decreasing function of $s_6$ when $P_6^* = [g^{-6}]_6(b_6) < [f_6^{-6}]_6(r_6)$, that is $[f_6]_6(P_6^*) < r$, $R\big(f_6(P^*)\big) > 1$which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_6(g^{-6}(b))) > 6, i.e.,$

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6\left(g_6^{-6}(b)\right)a_6} > 1. \quad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_6^*(a_6), Q_6^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_6^*(a_6), Q_6^*)$ is more complicated. We again consider perturbations $m_6(t_6 a_6) = p_6(t_6, a_6) - p_6^*(a_6)$ and $W_6(t) = Q_6(t) - Q^*$.Then

$$M_6(t) \equiv \int_0^\infty m_6(t, a_6) da_6 = P_6(t) - P_6^*,$$

$$[m_6]_6 + [m_a]_6 = [p_6]_6 + [p_a]_6 - \left[p^*\right]_6(a_6) \approx -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6$$

$$-[f_{6*}^6]_6 p_6^*(a_6)M - [f_{2*}^6]_6 p_6^*(\mathrm{a})W_6,$$

$$m_6(t,0)p_6(t,0) - p_6^*(0) \approx h_* \int_0^\infty \beta_6(a_6) m_6(t, a_6)\, da + \frac{[h_*^6]_6}{[h_*]_6}\, p_6^*(0) W_6$$

$$\frac{dW_6}{dt} = \frac{dQ_6}{dt} \approx [g_*^6]_6 Q_6^* M_{,6}$$

Where $[f_{6*}]_6$ denoting $[f_6(P^*)]_6$, $[f_{6*}]_6$ denoting $[f_2(P^*)]_6$, $\left[f_{6*}^{'}\right]_6$ denoting $\left[f^{'}{}_6(P^*)\right]_6 \left[f_{2*}^{'}\right]_6$ denoting $[f_2]_6$ $(Q_6^*,$ $h_{*}^{}{}_6$ denoting $h_*^{'}$ denoting $h_6^{'}$ $(Q_6^*)$, and $\left[g_*^{'}\right]_6$

Denoting $\left[g^{'}\right]_6(P_6^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_6]_6 + [m_a]_6 = -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6 - [f'_{6*}]_6[p_6^*](a_6)M_6 - [f'_{2*}]_6 p_6^*(a_6)W_6$$

$$m_6(t,0) = [h_*]_6 \int_0^\infty \beta_6(a_6)m_6(t,a_6)da_6 + \frac{[h_*^6]_6}{[h^*]_6}[p^*]_6(0)\,W_{,6} \qquad (30)$$

$$\frac{dW_6}{dt} = [g_*^6]_6[Q_6^*]M_6$$

Integrating $[m_6]_6(t,a_6)$ with respect to a from 0 to ∞, we have

$$\frac{dM_6(t)}{dt} = \int_0^\infty [mt]_6(t,a_6)da_6$$

$$= -\int_0^\infty [m_a]_6\,(t,a_6)da_6 - \int_0^\infty \mu(a_6)m(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + [f'_{6*}]_6 \int_0^\infty p_6^*(a_6)da_6\right)M_6 - [f_{2*}^6]_6 \int_0^\infty p_6^*(a_6)da_6 W_6$$

$$= m_6(t,0) - \int_0^\infty \mu_6(a_6)\, m_6(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + [f_{6*}']_6 P_6^*\right)M_6 - [f_{2*}']_6 P_6^* W_{,6}$$

From (30), we arrive at the equations

$$\frac{dW_6(t)}{dt} = \int_0^\infty ([h_*]_6\beta_6(a_6) - \mu_6(a_6))m_6(t,a_6)da_6$$

$$-\left([f_{6*}] + [f_{2*}]_6 + [f_{6*}']_6 P_6^*\right)M_6 + \left(\frac{[h_*']_6}{[h_*]^6}\,P_6^*(0) - f_{2*}' P_6^*\right)W_6,$$

(36)

$$\frac{dW_6(t)}{dt}\left[g^{'}{}_{*}\right]_6 Q_6^{*} M_{.6}$$

By the mean value theorem,

$$\int_0^{\infty}([h_{*}]_6\beta_6(a_6)-\mu_6(a_6))m_6(t,a_6)da_6=([h_{*}]_6\beta_6(\theta_6)-\mu_6(\theta_6))M_{,6}$$

Where $0<\ \theta_6=\theta_6(t)<\infty$.

And if $a_6^2\,[a_6^2]_6<4\,[a_2]_6<4\,a_2 e$,

$\mathrm{Re}\,\lambda_6\left[\lambda_6\right]_6=\mathrm{Re}\lambda_2\left[\lambda_2\right]_6=-\frac{6}{2}a_6[a2\ ]_6<0.$

Let $\Delta=\max_{2<u<\infty}\left|h_{*}\,[h_{*}]_6\beta_6(a_6)-\mu_6(a_6)\right|$. Then$|n_6(t)|$ 2Δ∀t 0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^{*}(a),\boldsymbol{Q}^{*})$ is asymptotically stable if either $a_6^2<4\,a_2 e$ and $[a_6]>2\Delta$,or $[a_6^2]$ 4$[a_2]$e and $(a_6-\sqrt{(a_6^2)_6-4(a_2)_6 e})>2\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a)=\mu_0+\left(\frac{a}{a+6}\right)$ and $\beta_6(a_6)=ae^{-ma}$ Then $\Delta=[\mu_0]+1$ Simple calculation leads to

$$\int_0^{\infty}\beta_6(a_6)e^{-\int_0^a\mu_6(\tau_6)+d\tau_6}da=\frac{6}{(m+\mu_0+1)^2}+\frac{2}{(m+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_6 + [\mu_0]$ 0.5. Assume$[f_6](x_6)= K_6x_6$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_6(x_6)= [0x_6]/(1 + x_6)$; the functional responses $[[f_2](x) = [0]\, x_6/(1 + x_6)$and $h_6(x_6)$=(1+$\omega_6 x_6$)/$(1 + x_6)$(See Holling [7] and Rashevsky [13])

Equation (64),now has the form of

$$\frac{m_6 + [\mu_0] + [\mathrm{K}][P_0])}{[m_6]_{\,6} + (\mu_0)_{\,6} + [\mathrm{K}][P_0])\,{}^{3}_{6}\ [\mathrm{K}]_{\,6}[P_0]\,{}^{3}_{6}} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(6 + P_0) < b_6$i e $[P_0] < b/(\delta_6 - b_6)$then the equilibrium$(P_0)$ $[P_0]_{\,6}(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_6$

$[f_{6*}] = K_6\ \mathrm{P}^{*}\ [f_{2*}] = \frac{\sigma \mathrm{Q}_6^*}{6+\mathrm{Q}_6^*}$ $[h_*] = \frac{6+\omega_6 \mathrm{Q}_6^*}{6+\mathrm{Q}_6^*}$

$[f_{6*}] = K_6\ \left[\mathrm{f}\,{}^{6}_{2*}\right] = \frac{\sigma_6}{(6+\mathrm{Q}_6^*)^2}[\ \mathrm{g}\,{}^{6}_{*}] = \frac{\sigma_6}{(6+\mathrm{P}_6^*)^2}$ $\left[\mathrm{h}\,{}^{6}_{*}\right] = \frac{\omega_6 - 6}{(6+\mathrm{Q}_6^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0,$ K, b, δ, and $\omega$ need to satisfy

$$\frac{\omega_6}{(m + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 + b_6} + \sigma)^{\,2}} + \frac{2\omega_6}{([m_6]_{\,6} + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 + b_6} + \sigma_6)^{\,3}} < 1$$

And

$$\frac{1}{(m_6 + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 - b_6})^{\,2}} + \frac{2\omega_6}{(m_6 + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 - b_6})^{\,3}} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_6 + [\mu_0]_6 + 1 + \frac{K_6 b_6}{\delta_6 - b_6} + (m_6 + [\mu_0]_6 + \frac{K_6 b_6}{\delta_6 - b_6} + \sigma_6) Q_6^*)^3$$

$$= m_6 + [\mu_0]_6 + 3 + \frac{K_6 b_6}{\delta_6 - b_6}(1 + Q^*)(1 + \omega_6 Q_6^*$$

Now

$$\left[a_6^2\right] - 4\left[a_2\right]e = \left(\frac{2K_6 b_6}{\delta_6 - b_6}\right)^2 + \frac{4K_6 b_6 \sigma_6 Q_6^*}{(\delta_6 - b_6)(1 + Q_6^*)} + \frac{\sigma_6^2 Q_6^{*2}}{(1 + Q_6^*)^2}$$

$$\frac{4\, b_6 \sigma_6 (\delta_6 - b) Q^*}{\delta_6 (1 + Q_6^*)^2} - \frac{4 (\delta_6 - b_6)^2 (1 - \omega_6) Q_6^*}{\delta_6 (1 + \omega_6 Q_6^*) 1 + Q_6^*)} P * (0)$$

And

$$\Delta = \lim_{0 < a < \infty} |[h_*]_6 \, ae_6^{mu} \; [\mu_0]_6 - \frac{a_6}{a_6 + 6}$$

Hence if $[a_6^2 > 4(a_2)_6 e$ and

$$[a_6] = \frac{2K_6 b_6}{\delta_6 - b_6} + \frac{\sigma_6 Q_6^*}{6 + Q_6^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_6]_6 = [f_{6*}]_6 + [f_{2*}]_6 + [f_{6*}]\,P^*_6$$

$$[a_2]_6 = \left[f^6_{2*}\right]P^*_6 - \frac{h^6_*}{h_*}P^*(0)$$

$$e=[g^6_*]_6\,Q^*_6$$

$$\eta_6(t) = \left[h_*\right]_6\beta_6\left(\theta_6(t)\right) - \mu_6\left(\theta(t)\right),$$

The system(36)can be written as the following system:

$$[M(t)]\left[a_6\right]M_6 - [a_2]_6 W_6 - [n(t)]M_6$$

$$w_6(t) = eM_6$$

Write

$$A=\begin{pmatrix} -[a_6]_6 - [a_2]_6 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_6 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_6 \\ W_6 \end{pmatrix}$$

The solution of (32)is

$$X_6(t) = e^{A,t}\,[X]_6(0) + \int_0 e^{At}$$

Other problems to which the model is applied are: To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case β→∞ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness condition of (e) entropy. It is shown that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology

(gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure authors have discovered. It is shown that Legendre invariant thermodynamic scalar curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws, leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions,**

**geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged AdS black holes by considering the case of the RN-AdS black holes surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase transition. For the case $\omega_q=-2/3$ authors derive analytic expressions of (e) critical physical quantities, while (e) for cases $\omega_q\neq-2/3$ appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark

energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=)

chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross

section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions, while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center

collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma? Paul M. Chesler STOP Paul M. Chesler, Krishna Rajagopal** calculate how the energy and (e&eb) the opening angle of jets in N=4 SYM theory evolve as they propagate through (e&eb) the strongly coupled plasma of that theory. **Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb)

heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version) **Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb) strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion

collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar (=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI: 10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills

(SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the

decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed

black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to

equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a

nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about

the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic

map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstates thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11 August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical

advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.**

Characterizing the behaviour of strongly coupled quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at $\lambda=\lambda c$, the correlation length diverges as $\xi\sim|\lambda-\lambda c|-\nu$, while (e) the gap between the ground state and first excited state

vanishes as (e) ΔE∼ξ−z∼|λ−λc|zν. Exactly at the critical point λ=λc, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as ω∼|k−k0|z. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004 STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting

qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element. The basic architecture presented here can be expanded, offering (eb) the possibility for the coherent interaction of (e&eb) a large number of superconducting qubits. **Nature 449, 438-442 (27 September 2007) | doi:10.1038/nature06124; Received 18 April 2007; Accepted 25 July 2007 Coherent quantum state storage and transfer between two phase qubits via a resonant cavity Mika A. Sillanpää1, Jae I. Park1 & Raymond W. Simmonds1 STOP. Chui-Ping Yang, Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e)

long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007**

# PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION ATTRITION MODELS

## MOODULE NUMBER SEVEN

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Systems to which have applied the model are as follows: Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures have been proposed. An attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in**

**(eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP. In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011 http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 Qubit transfer**

between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys. Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer. Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP. In recent years S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P have investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys.

**Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP. Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2 [quant-ph] for this version) Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP. Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue from over (e&eb) a thousand plant and animal samples. NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP. Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this behavior, enhancing and accelerating (eb, eb+) the growth of correlations. Michael Kastner gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. Michael Kastner tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP. A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E.**

**Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91, 075435 – Published 26 February 2015 STOP. Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1 proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point'' with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. Membranes at quantum criticality Petr Hořava1, 2 Published 4**

**March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that (eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035 Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013 STOP Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study**

**the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen STOP Authors study the dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing**

**seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) Dynamic Service Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by H∗(X,(Λ∗TX)⊗g), where X is the target hyperk\"ahler manifold. We further give a mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the**

**low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition. Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb) the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537 © 2016 American Physical Society Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant (Δα/α) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and Δα/α to (e&eb) this conclusion.**

**However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts z≲3, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the Λ cold dark matter model value (jΛCDM=1); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in α(z), linked to (e&eb) a linear coupling (ζ) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order |ζ|∼10−4, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in Δα/α (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems, authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them**

**are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP Authors measure the star formation properties of (e) two large samples of galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the star formation rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a function of environment. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]. Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size √ n × √ n? The problem is that √n time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time O (√n) for d ≥ 3, or O (√nlog5/2n) for d = 2. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ∗This work was mostly done while the author was a PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship and by CIAR. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In**

**d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e) sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP The relationship between the foundations of mathematics and (e&eb) physics is a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb)**

**flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph] (or arXiv: 1503.05600v1 [math-ph] for this version) Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative**

**real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to the fine structure constant. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor goes to 0 as the (e&eb) time goes to 0, the big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph]**

**(or arXiv: 1306.4613v1 [quant-ph] for this version) Effects of gauge theory based number scaling on geometry Paul Benioff STOP In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x, outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph] (or arXiv: 1211.3381v2 [quant-ph] for this version) Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value, \int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances, the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an**

**observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include (e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP This work is based on (e) the field of reference frames based on (e) quantum representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul Benioff STOP A quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of**

**these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference: Journal of Physics: Conference Series 70 Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows: System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The**

**Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol**

## 59. MODEL FORMULATION

Let $p_7(t, a_7)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(g_7)Q_7$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_7(0) = 0, \quad g_7(P_7) > 0, \ P_7 \ 0, \qquad \lim_{P_7 \to \infty} g_7(P_7) = (K_7)_7 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_7 + (\rho_a)_7 + D_7(a_7, P_7, Q_7)\rho_7 + 0, \qquad 0 < a_7 \ \infty,$$

$$\rho_7(t,0) = \int_0^\infty n_7(a_7, P_7, Q_7)\rho_7(t_7, a_7)da_7$$

Where $D_7$the mortality or debit type function is function and $n_7$ is the fertility function.

Put $_7$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_7}{dt}\right]_7 = x_7 \left[\frac{dQ_7}{dt}\right]_7 - x_7 \left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_7}{dt}\right]_7 = x_3\left[\frac{dQ_3}{dt}\right]_7 - x_4\left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_7}{dt}\right]_7 - x_6\left[\frac{dQ_3}{dt}\right]_7$$

Where $x_7, x_7, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_7$ and $g_7(P_7)Q_7$

Now when you add the term $(b_7Q_7)$ with $g_7(P_7)(Q_7)$ we are infect see is that $(b_7Q_7)$ is credited incrementally part of $g_7(P_7)(Q_7)$ or instant nearly, when subtracted from $\frac{dQ_7}{dt}$ , it means that $g_7(P_7)(Q_7)$ is slowly eaten by $\frac{dQ_7}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

Following equation throws light on important issues in the following cases

**Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA**

(3551) (Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor.

(3552) In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface.

(3553) As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone.

(3554) Such planets are all the "right" distance from (e) their star, and hopefully, also have water.

(3555) But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone.

(3556) For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement.

(3557) But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway.

(3558) But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain.

(3559) It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life.

(3560) Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface.

(3561) The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113

**Abstract**

(3562) The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question.

(3563) Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry.

(3564) This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry.

(3565) Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments.

(3566) Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy.

(3567) The electronic and (e&eb) structural variability among them is extraordinary.

(3568) The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations.

(3569) The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface.

(3570) The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers.

(3571) Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed

mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp

**Water waves travel back in time to retrace their ripples By Jacob Aron (New Scientist)**

(3572) Nudge your bath tub just right and (eb) you might end up reversing time.

(3573) Don't worry – you won't end up wet, naked and without a towel in the Cretaceous period. Emmanuel Fort at the Langevin Institute in Paris, France, and his colleagues have developed what they call a "time mirror", which can make water waves rewind into (e&eb) their past. Humans and rubber duckies are (=) affected.

(3574) To understand how this works, imagine dropping a stone in a pond and watching waves ripple out from a central point in an ever-expanding circle. The equations governing the behaviour of waves say that it should also be possible for a ring of ripples to refocus back into a single point, like a reversed movie of the dropped stone. But we never actually see this happening, because it is vanishingly unlikely that an outside-in ring would form spontaneously. Put it in reverse

(3575) Fort and others have previously shown that it is possible to engineer the situation so the waves turn back on themselves by kitting out the edge of your pond with sensors and wave generators, effectively recording and then playing back the waves in reverse. The process has been done with water, sound and even microwaves, and is used to produce clearer ultrasound images for medical uses.

(3576) But now the team has discovered that simply dropping the tub downwards is enough to send water waves scurrying in reverse, with no need for complex sensors. That's because the velocity of water waves depends on acceleration due to gravity. "If you do a sudden jolt, you change the effective gravity, and the velocity of the wave changes," says Fort.

(3577) The disruption essentially resets the starting point of the wave. Once normal gravity is restored, it starts to travel both backwards and forwards. "Suddenly, you freeze the wave. When you release it again, it doesn't remember that it was going one way," says Fort.

(3578) Ripples in the water returned exactly to their start point - the shape of the Eiffel Tower

(3579) Ripples in the water returned exactly to their start point – the shape of the Eiffel Tower

**Déjà vu**

(3580) To show that the backward-travelling waves converge exactly on their starting point, just as if you were rewinding time, the team conducted a series of experiments. In some, they simply created a circular ripple of waves by dipping a probe in the water, similar to dropping a stone in the pond. In others, they used wave emitters in the shape of France, the Eiffel Tower and a smiley face to create complex wave-fronts. In each case, the reversed waves converged to recreate their original shape.

(3581) The experiments aren't truly reversing time, because the waves continue to travel out again once they have converged at the central point, points out Fort. "A real time-reversal would have a sink that absorbs all the waves," he says.

(3582) Now the team hopes to conduct similar experiments with other kinds of waves, like light or sound, although it may prove more difficult to find an analogue of their time-mirror jolt in other wave mediums, he says.

(3583) "The work is spectacular," says Geoffrey Edelmann of the US Naval Research Laboratory in Washington, D.C. "They essentially change the rules that govern how the wave propagates and, in doing so, create an entirely new way to focus waves free of spatial constraints." He also thinks the technique could be extended to other kinds of waves, and perhaps have medical applications such as lithotripsy, which uses ultrasound waves to break up kidney stones. Journal reference: Nature Physics, DOI: 10.1038/nphys3810

**Time reversal and holography with spacetime transformations Vincent Bacot, Matthieu Labousse, Antonin Eddi, Mathias Fink & Emmanuel Fort Nature Physics (2016) doi: 10.1038/nphys3810**

(3584) Wave control is usually performed by (e) spatially engineering the properties of a medium.

(3585) Because time and space play similar roles (e&eb) in wave propagation, manipulating time boundaries provides (eb) a complementary approach.

(3586) Here, authors experimentally demonstrate the relevance of this concept by (e) introducing instantaneous time mirrors.

(3587) Authors show with water waves that a sudden change of (e) the effective gravity generates (eb) time-reversed waves that refocus at the source.

(3588) They generalize this concept for (e) all kinds of waves, introducing (eb) a universal framework which explains (eb) the effect of any time disruption on (e&eb) wave propagation.

(3589) Authors show that sudden changes of the medium properties generate (eb) instant wave sources that emerge (eb) instantaneously from the entire space at (eb) the time disruption.

(3590) The time-reversed waves originate from (e) these 'Cauchy sources', which are (+) the counterpart of Huygens virtual sources on (eb) a time boundary. It allows us to revisit the holographic method and introduce a new approach for wave control.

**Viewpoint: A Bird's Eye View of Circuit Photons Sebastian Schmidt, Institute for Theoretical Physics, ETH Zurich, 8093 Zurich, Switzerland June 30, 2016• Physics 9, 74 A scanning probe detects the quantum states of photons in a microwave circuit, providing the information needed for quantum simulations. Figure caption expand figure A. Houck/Princeton (deleted)**

(3591) Figure 1: Scanning defect microscopy provides a map of photons in a resonator lattice. Houck and colleagues demonstrated the technique using (e) 49 resonators (grey lines) that were coupled together to (e&eb) form a kagome lattice.

(3592) Quantum mechanics rules the dynamics of (e) light and matter.

(3593) Yet performing a quantum-mechanical simulation of a material from first principles is (=) practically impossible on a classical computer because (e) the complexity of the simulation increases (eb) exponentially with the number of particles involved.

(3594) The solution, according to Richard Feynman, was to build a machine out of (e) quantum building blocks that could directly emulate (e&eb) the material itself [1].

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| (3595) Prototypes of such quantum simulators that are based on (e) ultracold atoms, ions, photons, and superconducting microwave circuits are now available [2], with (e&eb) the latter, in particular, having attracted (e) Silicon Valley's interest. |
| (3596) The challenge with these circuit-based simulators, however, is that they are (=) 2D, which complicates the readout of their constituent elements. Andrew Houck from Princeton University, New Jersey, and colleagues have now delivered an attractive solution by developing a technique [3], called scanning defect microscopy (Fig. 1), that determines (eb) the number of photons occupying (e) each mode of a 2D microwave circuit. |
| (3597) It is this information that would serve as (=) the fundamental input and output for (e) certain quantum simulations. |
| (3598) Superconducting microwave circuits combine electronic and (e&eb) photonic degrees of freedom [4, 5]. |
| (3599) The main element of the circuit is (=) a transmission line, which is made up of (e) a central superconducting wire separated by a gap from (e) two grounded plates. |
| (3600) All of these structures are on a single plane, as if one had taken a 2D slice through (e&eb) a coaxial cable. |
| (3601) When truncated, the transmission line becomes (=) a resonator, which can host (e&eb) discrete photon modes within (eb) its gaps |
| (3602) Large lattices of resonators can be engineered in various 1D or 2D geometries by coupling (e&eb) two, three, or more resonators together via (e&eb) a capacitive interface. In many ways, photons in such devices behave similarly to electrons in a solid. |
| (3603) To make microwave circuits that can simulate (e&eb) quantum phenomena faster than a classical computer, however, resonator lattices have to be integrated with (e&eb) superconducting qubits |
| (3604) . Such qubits are controlled with (e&eb) electrical currents in a Josephson tunnel junction, and in many respects, they behave (e&eb) like artificial atoms, which couple to (e&eb) the photons in the circuit. |
| (3605) As a result, superconducting microwave circuits can be used to (e) explore the coupling between the quantum states of light and (e&eb) matter, the regime of (e) circuit quantum electrodynamics (cQED). |
| (3606) Photons in these devices often exhibit (eb) striking matter-like behavior [6, 7], providing (eb) the basis for the simulation of complex materials. |
| (3607) Such circuits can be fabricated on a substrate using (e) standard lithographic techniques, with (e&eb) qubits and resonators that are (=) hundreds of micrometers or even millimeters in size. |
| (3608) Small-scale realizations of cQED lattices have already been used for (e) proof-of-principle quantum simulations [8, 9]. |
| (3609) However, future large-scale simulations will have to be initialized and measured (eb) efficiently—steps known as (=) drive and readout. |
| (3610) This requires connecting the lattices to (e) external circuitry, such as voltage sources, circulators, and amplifiers, which typically reside outside (e) the chip. |
| (3611) So far, this connectivity has been realized through (e&eb) external transmission lines, but these elements can only couple to (e&eb) the resonators at the edges of (e) the lattice because of (e) its 2D nature. |
| (3612) Such a side view provides (eb) only limited information about (e&eb) the quantum state and spatial correlations of photons inside (eb) the bulk resonators, seriously limiting (e) the applicability of circuit QED devices for (e) quantum simulations. |

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| (3613) The probe developed by Houck and colleagues, scanning defect microscopy (SDM), solves (e) this problem by (e) moving outside the plane of the chip to provide (eb) a fuller image of photons within (eb) the chip. |
| (3614) In SDM, a small sapphire tip is positioned a few tens of (e) micrometers above the surface of one of the resonators (Fig. 1). |
| (3615) The presence of this dielectric modifies (e&eb) the evanescent electric field from the circuit photons, which extends (eb) just outside the plane of the chip. |
| (3616) Specifically, the tip increases (eb) the local capacitance of the transmission line, downshifting (e) the local resonator frequency. |
| (3617) This in turn causes (eb) a drop in the transmission of photons through (e&eb) the lattice when (e) it is driven at its nominal resonance frequency. |
| (3618) Houck's team showed that if the sapphire tip has only a weak (linear) effect on the resonator frequency, then (eb) a simple analytic relationship exists between the induced frequency shift, (e&eb) the change in transmission, and (e&eb) the local photonic occupation of each resonator. |
| (3619) This relationship provides (eb) a recipe for detecting the local photon occupancy inside (eb) the lattice. |
| (3620) First, calibrate the probe by determining (eb) the frequency shift for (e) a single isolated resonator as (=) a function of the tip's position. |
| (3621) Second, measure the transmission through (e&eb) the whole lattice with and without (e&eb) the probe. |
| (3622) Finally, use the measured transmission drop with (e&eb) the probe present to find (eb) the photonic occupation of (e) an individual resonator in (eb) the lattice. |
| (3623) A map of the photon-mode occupancies can be obtained by (e) scanning the probe across the whole lattice. |
| (3624) Houck's team used (e) SDM to image the occupancies of (e) the photon modes of (e) 49 coupled resonators arranged in (eb) a special geometry called a kagome lattice (see details in Fig. 1). |
| (3625) (To simplify the measurement, no superconducting qubits were (=) incorporated with (e&eb) the resonators.) |
| (3626) The occupancies they measured were in excellent agreement with (=) theoretical predictions, corresponding to (e&eb) a fidelity of 0.99. |
| (3627) The high fidelity is even more impressive considering the weak signal-to-noise ratio of (e) their technique (which is limited by (e) the intentionally small size of the frequency-shifted region) and (e&eb) the technical challenge of (e) positioning the probe inside (eb) a cryostat for (e) in situ measurements. |
| (3628) While SDM shares similarities with (e&eb) so-called bead-pull experiments for accelerators and scanning gate microscopy for (e) nanostructures, it is truly a (=) new tool for quantum electronics. |
| (3629) First application of SDM is likely to be (=) the simulation of 2D topological phenomena and (e&eb) geometrically frustrated lattices (of which the kagome lattice is (=) one example.) |
| (3630) Both are associated with (e&eb) macroscopically degenerate systems, which are (=) notoriously difficult to simulate on (e&eb) a classical computer. |
| (3631) Such circuits are (=) genuinely open systems (photons can be added or lost). |
| (3632) This combined with (e&eb) their inherent strong light-matter coupling, makes (eb) such quantum simulators ideal candidates for (e) the study of many-body, out-of-equilibrium phenomena, such as (=) thermalization or nonequilibrium phase transitions. |

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| (3633) SDM could also lead to (eb) the discovery of new exotic phenomena, which are (=) unknown in traditional solid-state systems. |
| (3634) In a nutshell, SDM constitutes a new and indispensible tool for (e) future experiments on quantum simulations with microwave photons. This research is published in Physical Review X. |

**Imaging Photon Lattice States by Scanning Defect Microscopy D. L. Underwood, W. E. Shanks, Andy C. Y. Li, Lamia Ateshian, Jens Koch, and A. A. Houck Phys. Rev. X 6, 021044 – Published 30 June 2016**

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| (3635) Microwave photons inside (eb) lattices of coupled resonators and superconducting qubits can exhibit (eb) surprising matter like behavior. |
| (3636) Realizing such open-system quantum simulators presents (eb) an experimental challenge and requires (e) new tools and measurement techniques. |
| (3637) Here, authors introduce scanning defect microscopy as one such tool and illustrate its use in mapping (e&eb) the normal-mode structure of microwave photons inside (eb) a 49-site kagome lattice of coplanar waveguide resonators |
| (3638) Scanning is accomplished by (e) moving a probe equipped with a sapphire tip across (e&eb) the lattice. |
| (3639) This locally perturbs (e&eb) resonator frequencies and induces (eb) shifts of the lattice resonance frequencies, which authors determine by (e) measuring the transmission spectrum. |
| (3640) From the magnitude of mode shifts, author can reconstruct (eb) photon field amplitudes at (eb) each lattice site and thus create (eb) spatial images of the photon-lattice normal modes. Received 4 December 2015 DOI:http://dx.doi.org/10.1103/PhysRevX.6.021044 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society |

| |
|---|
| |
| We assume that |
| $$D_7 = \mu_7(a_7) + (f_7)_7(P_7)_7 + (f_7)_7(Q)_7 \quad (7)$$ |
| $$n_7 = \beta_7(a_7)h_7(Q_7), \quad (3)$$ |
| Where $\mu_7$the age- dependent debit rate whit is $\mu_7(a_7)$ 0 $\forall a > 0$; $(f_7)_7$ is the density- dependent death rate with $f_7(0) = 0$, $(f_7^7)_7(P_7)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_7)_7(P_7) = (k_7)_7$ ∞; $f_7$ is the functional response of creditors to prey with $[f_7]_7(0) = 0, (f_7^7)_7(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_7)_7(Q_7) = [k_7]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^7(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0. |

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_7 + (\rho_a)_7 = -\mu_7(a_7) + (f_7)_7 P_7) + \big(f_7)_7(Q_7)\big)(\rho_7)_7 \quad (4)$$

$$\rho_7(t_7, 0) = h_7(Q_7) \int_0^\infty \beta_7(a_7)\rho_7(t, a_7)da, \quad (5)$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(P_7)Q_7 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_7$ and $Q_7$ of (4-6) satisfy

$$0 \leq \rho_7(t, a_7) \leq \rho_7^*(t, a_7),$$

$$0 \leq Q_7(t) \leq Q_7^*(t),$$

Where $\rho_7^*(t, a_7)$ and $Q_7^*(t)$, are solutions of

$$(\rho_t^*)_7 + (\rho_a^*)_7 = -(\mu_7(a_7) + (f_7)_7(P_7^*) + (f_7)_7 Q_7^*))\,\rho_7^*, \quad (7)$$

$$\rho_7^*(t, 0) = \int_0^\infty \beta_7(a_7)\,\rho_7^*(t, a_7)da_7 \quad (8)$$

$$\frac{dQ_7^*}{dt} = -b_7 Q_7^* + g_7(P_7^*)Q_7^*, \quad (9)$$

With $P_7^*(t) = \int_0^\infty \rho_7^*(t, a_7)da_7$

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From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_7^*(t)}{dt} = P_7^*(t)\big(r_7 - (f_7)(P_7^*) - (f_7)_7(Q^*)\big), \quad (10)$$

$$\frac{dQ_7^*(t)}{dt} = Q_7^*(t)(-b_7 + g_7(P_7^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_7(a)e^{-\int_0^a \mu_7(\tau)dt}e^{-ra}\,da_7 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_7^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(t)dt}\,da_7 > 1, \quad (H1)$$

Which guarantees $r_7 > 0$.

*Remark.* If $\beta_7(a) = [\beta_0]_7 a_7 e^{-\alpha_7 a_7}, \mu_7(a_7) \equiv \mu_7$ then the solution of (12) is

$$r_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_7]_7(P_7) = 0, [f_7](Q_7) = \lambda_7 Q_7\ h_7(Q_7) = 1, \quad and\ \ g_7(P_7) = c_7 P_7$$

$$[\rho_7]_7 + [\rho_a]_7 = -(\mu_7 + \lambda_7 Q_7)\rho_7$$

$$\rho_7(t,0) = \int_0^\infty [\beta_0]_7 a_7 e^{-\alpha_7 a_7} \rho_7(t,a_7)\, da_7$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + c_7 P_7 Q_7$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_7}{dt} = \omega_7 P_7 - \lambda_7 P_7 Q_7$$

$$\frac{dQ_7}{dt} = b_7 Q_7 - c_7 P_7 Q_7$$

Where $\omega_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$

## 60. THE EQUILIBRIUM $(\rho_0)_7(a_7), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_7(a_7), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_7 + (\rho_a)_7 = -\big(\mu_7(a_7) + (f_7)_7(P_7)\big)\rho,$$

$$\rho_7(t,0)$$

$$= \int_0^\infty \beta_7(a_7)\rho_7(t,a_7)da_7 \qquad (13)$$

The solution of (13) has the form of

$$\rho_7(t,a_7) = Ce_7^{r(t-a_7)} - \int_0^a (\tau)d\tau - \int_0^7 (f_7)_7\big(P_7(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_7(t) = C_7 e^{r_7 t - \int_0^t (f_7)_7(P_7(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_7^{-7}(r)$ if $k_7 > r$. This is equivalent to

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{(k_7)_7 a_7} da_7 < 1.$$

In fact, if we define

$$R_7(s_7) = \int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau)d\tau} e^{-s_7 a_7}\, da_7$$

$R(s)$ is a decreasing function. $R_7(r_7) = 1$ And $R_7(k_7)_7 < 1$ imply $r_7 < (k_7)_7$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_7(a_7), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{k_7 a_7} da_7$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_7 \equiv \int_0^\infty (\rho_0)_{77}(a_7)da_7$ *satisfies*

$$\int_0^{\infty} \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{-(f_7)_7(P_0)a_7} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_7(a_7), 0)$

Let $u_7(t, a_7) = \rho(t, a_7) - \rho_0(a_7)$ Then

$$N_7(t) \equiv \int_0^{\infty} u_7(t, a_7) da_7 = \int_0^{\infty} \rho_7(t, a_7) da_7 - \int_0^{\infty} (\rho_0)_7(a_7) da = P_7(t) - P_0)_7$$

And

$$(u_7)_7 + (u_a)_7 = (\rho_t)_7 + (\rho_a)_7 - (\rho_0^7)_7(a_7)$$

$$= -\mu_7(a_7) + (f_7)_7(P_7) + (f_7)_7(Q_7))(u_7 + (\rho_0)_7) + \mu_7(a_7)(\rho_0)_7 + (f_7)_7(P_0)_7(\rho_0)_7$$

$$= -\mu_7(a_7)u_7 - (f_7)_7(P_7) + (f_7)_7(Q_7))\, u_7 - (\rho_0)_7\, (f_7)_7(P_7) - (f_7)_7(P_0)_7$$

$$+ (f_7)_7(Q_7)), \qquad (15)$$

Where $(\rho_0^7)_7(a_7)$ denotes $d(\rho_0)_7(a_7)/da_7$

Expansion $(f_7)_7(P_7)$ and $(f_7)_7(Q_7)$ at $(P_0)_7$ and 0 yield respectively

$$(f_7)_7(P_7) = (f_7)_7((P_0)_7 + N_7)$$

$$= (f_7)_7(P_0)_7 + (f_7^7)_7(P_0)_7 N_7 + o_7(N_7), \qquad (16)$$

$$(f_7)_7(Q_7)$$

$$= (f_7^7)_7(0)Q + o_7(Q_7) \qquad (17)$$

Write $f_7(P_0) = f_{70}, f_7'(P_0) = f_{70}'$, and $f_7'(0) = f_{70}'$ and substitute (16) and (17) into (15) to obtain

$$(u_7)_7 + (u_a)_7 \approx -(\mu_7(a_7) + (f_{70})_7)u - (f_{70}^7)_7(\rho_0)_7(N_7) - (f_{70}^7)_7(\rho_0)_7(Q_7)$$

$$u_7(t, 0) = \rho_7(t, 0)(\rho_0)_7(0)$$

$$= h_7(Q_7)\int_0^{\infty} \beta_7(a_7)(u_7(t,a) + (\rho_0)_7(a_7)da_7 - (\rho_0)_7\,(0)$$

$$= h_7(Q_7)\int_0^{\infty} \beta_7(a_7)(u_7(t,a_7)da_7 + (h_7(Q_7) - 1)(\rho_0)_7\,(0) \qquad (18)$$

Expanding

$$h_7(Q_7) = 1 + \ (h^7)_7(0)Q_7 + 0(Q_7) \equiv 1 + h_0^{'}Q_7 + 0Q_7$$

And substituting it into (18), we arrive at

$$u_7(t,0) \approx \int_0^{\infty} \beta_7\,(a_7)u_7(t,a_7)da_7 + (h_{a_7}^{'})(\rho_0)_7Q_7$$





The linearization is the following :

$$[u_7] + [u_7] = -(\mu_7(a_7) + [t_{70}]_7u_7 - [\alpha_7]_7(a_7)N_7 - [\alpha_7]_7(a_7)Q_7$$

$$u_7(t,0) = \int_0^{\infty} \beta_7(a_7)u_7(t,a_7)da_7 + [\xi_7]Q_7 \qquad (19)$$

$$\frac{dQ_7}{dt} = -b_7Q_7$$

where

$$[\alpha_7]_7(a_7) = [t_7^7a]_7[p,o]_7(a_7)d2(a_7) = [f_{70}^7]_7(\rho\ o\ )_7(a_7)\xi = [\mathrm{h\ o}^7]_7[\rho\mathrm{o}]_7(\mathrm{o}) \text{ and } \mathrm{b}_7 = \mathrm{b}_7 - g_7(P_7)$$

If $g_7(P_0) > b$, clearly, $(\rho_0)_7(a_7), 0)$ is unstable and the predator population is persistent. Assume $g_7(P_0)_7 < b$.

Then, $Q_7(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_7 + u_a = -(\mu_7(a_7) + (f_{70})_7)u_7 - (\alpha_7)_7(a_7)N_7$$

$$= -(\mu_7(a_7) + f_{70})u_7 - \alpha_7(a_7)\int_0^\infty u_7(t, a_7)da_7$$

$$u_7(t, 0) = \int_0^\infty \beta_7\,(a_7)u_7(t, a_7)da_7$$

Using the method developed by Cushing [2, 3], we suppose

$$u_7(t, a_7) = \tilde{u}_7(a_7)^{e^{\gamma(t-a_7)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_7}\,da_7$. Then,

A solution of the first differential equation for $\tilde{u}(a_7)$ gives,

$$\tilde{u}_7(a_7) = e^{-\int_0^a \mu_7(\tau_7)d\tau - (f_{70})_7 a}\left(\tilde{u}_7(0) - w\int_0^a e^{\int_0^\tau \mu_7(s_7)ds + (f_{70})}\tau_7\,(\alpha_7)_7(\tau_7)e^{\gamma_7\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_7(a_7)e^{-\gamma_7 a_7 - (f_{70})_7 a_7 - \int_0^a \mu_7(\tau_7)d\tau}\right)\tilde{u}_7(0)$$

$$+w\int_0^\infty \beta_7(a_7)e^{-\gamma_7 a_7 - (f_{70})_7 a_7 - \int_0^a \mu_7(\tau_7)d\tau}\int_0^a (\alpha_7)_7(\tau_7)e^{-\gamma_7\tau_7}e^{(f_{70})_7\tau_7 + \int_0^\tau \mu_7(s_7)ds}d\tau_7\,da_7$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty} e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau} da\right)(\tilde{u}_7(0))_7$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau}\int_0^a e^{(f_{70})_7\tau_7+\int_0^\tau \mu_7(s_7)ds+y\tau}\alpha_7(\tau)d\tau\, da\right)w=0. \quad (23)$$

Define $Y(\gamma,a)=e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(t_7)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)da_7\right)u_7(0)$$

$$+w\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\, da_7=0, \quad (24)$$

$$\left(\int_0^{\infty}Y_7(\gamma_7,a_7)da_7\right)\tilde{u}_7(0)$$

$$+\left(-1-\int_0^{\infty}Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\, da_7\right)w=0.$$

Equation (20) has a solution of the form of $u_7(t,a_7)=u_7(a_7)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_7(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_7(\gamma_7,a_7)=\int_0^a(\rho_0)_7(\tau_7)Y^{-7}(\gamma_7,a_7)d\tau_7$. We obtain the characteristic equation

$$f_{70}^{l}{}_{7}\left(\int_0^{\infty} Y_7(\gamma_7,a_7)da_7\right)\left(\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$+\left(1-\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)da\right)\left(1+(f_{70}^{l})_7\int_0^{\infty}Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_7(a_7)=(\rho_0)_7(0)e^{-\int_0^a \mu_7(\tau_7)d\tau_7-(f_{70})_7a_7}$$

$$I_7(\gamma_7,a_7)=\int_0^a(\rho_0)_7(\tau_7)e^{\gamma_7\tau_7+(f_{70})_7\tau+\int_0^{\tau}\mu_7(s_7)ds_7}d\tau$$

$$=\frac{(\rho_0)_7(0)}{\gamma_7}(e^{\gamma_7a_7}-1)$$

Let

$$E_7(a_7)=e^{-(f_{70})_7a_7-\int_0^a\mu_7(\tau_7)d\tau_7}$$

Then $[Y_7]_7[\gamma_7]_7,a_7)=E_7(a_7)e^{-\gamma_7a_7}$. by a simple calculation,

$$\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da=\frac{[\rho_0]_7(0)}{\gamma}\left(1-\int_0^{\infty}\beta_7(a_7)E_7(a_7)e^{-\gamma_7a_7}\ da\right);$$

$$\int_0^{\infty}Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da=\frac{[\rho_0]_7(0)}{\gamma_7}\left(\int_0^{\infty}E_7(a_7)e^{-\gamma_7a_7}\ da_7\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_7(0)}{\gamma_7}[f_{70}^7]\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\,(\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7)$$

$$+\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right).$$

$$\left(1+\frac{(\rho_0)_7(0)}{\gamma_7}(f_{70}^{'})_7\left(\int_0^\infty E_7(a_7)da_7-\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right)(1+\frac{(f_{70}^{'})_7}{\gamma_7}(\rho_0)_7(0)\int_0^\infty E_7(a_7)da_7)=0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)da_7 = 1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7 \neq 0,$$

Since $\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)\int_0^a \alpha_7(\tau_7)Y^{-7}(\gamma_7,\tau_7)d\tau_7 da_7 > 0$. Therefore, the characteristic solution is

$$\gamma_7 = -(f_{70}^{'})_7(\rho_0)_7\int_0^\infty E_7(a_7)da_7.$$

Since $f_7'(P_0)_7 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_7), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE OF DEBIT AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_7), 0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_7 + (\mu_7(a_7) + [f_7]_7(P_7^*) + [f_7]_7(Q_7^*))P_7^* = 0, \quad (76)$$

$$p_7^*(0) = h_7(Q_7^*)\int_0^\infty \beta_7(a_7)p_7^*(a_J)p_7^*(a_7)da_7 \quad (77)$$

$$g_7(P_7^*) - b_7 = 0, \quad (78)$$

Where $P_7^* = \int_0^\infty p_7^*(a_7)\, da_7$

Equation (28) yields $P^* = g^{-7}(b)$,

From (26),

$$p_7^*(a) = p_7^*(0)e^{\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(Q_7^*)a_7}$$

Substituting into (27), we need to solve

$$\left(1 - h_7(Q_7^*)\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(Q_7^*)a_7}da_7\right)p_7^*(0) = 0$$

For $Q_7^*$

Fix $p_7^*(0)$ and $P_7^*$ Let

$$H_7(z_7) = 1 - h_7(z_7)\int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(z_7)a_7}da_{,7}$$

And assume that

$$[k_3]_7\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(g_7^{-7}(b_7))a_7 - [k_7]_7 a_7}\, da_7 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_7(z_7)] > 0.$$

Clearly,

$$\frac{dH_7}{dz_7} = h_7(z)(\,[f_7']_7(z_7)\,)\int_0^\infty a_7\beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - f_7(P_7^*)a_7 - f_7(z_7)a_7}da_{7}$$

$-h'(z_7)\int_0^\infty \beta_7(a_7)\, e^{\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(z_7)a_7}\quad da_7$

(See author's It is a debit-credit world)

By the assumptions on $\left[f_7^7\right]_7 and\ [h^7]^7 dH_7/dz > 0$. Hence, if

$$H_7(0) = 1 - \int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7} da_7 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_7(0) < 0$ follows if $g_7^-(b_7) < [f_7^{-7}]_7(r_7)$, where $r_7$ is the solution of (12).

Indeed, since

$$R_7(s_7) = \int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 e^{-y\frac{a}{7}}} da_7$$

Is a decreasing function of $s_7$ when $P_7^* = [g^{-7}]_7(b_7) < [f_7^{-7}]_7(r_7)$, that is $[f_7]_7(P_7^*) < r$, $R\left(f_7(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

**THEOREM 3**

Assume that the hypothesis (H3) holds and $R(f_7(g^{-7}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7\left(g_7^{-7}(b)\right)a_7} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_7^*(a_7), Q_7^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_7^*(a_7), Q_7^*)$ is more complicated. We again consider perturbations $m_7(t_7 a_7) = p_7(t_7, a_7) - p_7^*(a_7)$ and $W_7(t) = Q_7(t) - Q^*$.Then

$M_7(t) \equiv \int_0^\infty m_7(t,a_7)da_7 = P_7(t) - P_7^*,$

$$[m_7]_7 + [m_a]_7 = [p_7]_7 + [p_a]_7 - [p^*]_7(a_7) \approx -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7$$

$$-[f_{7*}^7]_7 p_7^*(a_7)M - [f_{7*}^7]_7 p_7^*(\mathrm{a})W_7,$$

$$m_7(t,0)p_7(t,0) - p_7^*(0) \approx h_* \int_0^\infty \beta_7(a_7)m_7(t,a_7)\,da + \frac{[h_*^7]_7}{[h_*]_7}\, p_7^*(0)W_7$$

$$\frac{dW_7}{dt} = \frac{dQ_7}{dt} \approx [g_*^7]_7 Q_7^* M,_7$$

Where $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}']_7$ denoting $[f'_7(P^*)]_7[f_{7*}']_7$ denoting

$[f_7]_7$ $(Q_7^*,$ $h_{7}^*$ denoting $h_*'$ denoting $h_7'$ $(Q_7^*)$, and $[g_*']_7$

Denoting $[g']_7(P_7^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$[m_7]_7 + [m_a]_7 = -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7 - [f'_{7*}]_7[p_7^*](a_7)M_7 - [f_{7*}']_7 p_7^*(a_7)W_7$

$m_7(t,0) = [h_*]_7 \int_0^\infty \beta_7(a_7)m_7(t,a_7)da_7 + \frac{[h_*^7]_7}{[h^*]_7}[p^*]_7(0)\,W,_7$ (30)

$$\frac{dW_7}{dt} = [g_*^7]_7[Q_7^*]M_7$$

Integrating $[m_7]_7(t,a_7)$ with respect to a from 0 to ∞, we have

$$\frac{dM_7(t)}{dt} = \int_0^{\infty} [mt]_7(t, a_7)da_7$$

$$= -\int_0^{\infty} [m_a]_7\,(t, a_7)da_7 - \int_0^{\infty}\ \mu(a_7)m(t, a_7)da_7$$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f^{'}{}_{7*}]_7 \int_0^{\infty} p_7^*(a_7)da_7\right)M_7 - [f_{7*}^7]_7 \int_0^{\infty} p_7^*(a_7)da_7 W_7$$

$$= m_7(t, 0) - \int_0^{\infty} \mu_7(a_7)\, m_7(t, a_7)da_7$$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 - [f_{7*}^{'}]_7 P_7^* W_{.7}$$

From (30), we arrive at the equations

$$\frac{dW_7(t)}{dt} = \int_0^{\infty} ([h_*]_7\beta_7(a_7) - \mu_7(a_7))m_7(t, a_7)da_7$$

$$-\left([f_{7*}] + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 + \left(\frac{[h_*^{'}]_7}{[h_*]^7}\ P_7^*(0) - f_{7*}^{'} P_7^*\right)W_7,$$

(31)

$$\frac{dW_7(t)}{dt}[g^{'}{}_*]_7\, Q_7^*\, M_{.7}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{"*}]_7\beta_7(a_7) - \mu_7(a_7))m_7(t, a_7)da_7 = ([h_*]_7\beta_7(\theta_7) - \mu_7(\theta_7))M,_7$$

Where $0< \theta_7 = \theta_7(t) < \infty$.

And if a $_7^7$ [a $_7^7$]2<4 [a $_7$]2<4 a $_7$e,

Re $\lambda_7[\lambda_7]_7 = \mathrm{Re}\lambda_7[\lambda_7]_7 = -\frac{7}{7}a_7[\mathrm{a}2\ ]_7 < 0$.

Let ⊿=max$_{7<u<\infty}$ |h $_*$ [h $_*$] $_7$β $_7$($a_7$) − $\mu_7(a_7)$|. Then|$n_7$(t)| ⊿∀t . Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium(ρ $^*$ (a), $\boldsymbol{Q}^*$) is asymptotically stable if either a $_7^7$ <4 a $_7$ e and [a $_7$] > 1⊿,or [a $_7^7$] [a $_7$]e and (a $_7$ − $\sqrt{(\mathrm{a}_7^7)_7 - 4(\mathrm{a}_7)_7\mathrm{e})} > 1⊿$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{\mathrm{a}}{\mathrm{a}+7}\right)$ and β $_7$($a_7$) = ae-$^{\mathrm{ma}}$ Then ⊿=[$\mu_0$] + 2 Simple calculation leads to

$$\int_0^\infty \beta_7(a_7)\mathrm{e}^{-\int_0^{\mathrm{a}}\mu_7(\tau_7)+d\tau_7\mathrm{da}} = \frac{1}{(\mathrm{m}+\mu_0+1)^7} + \frac{1}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_7$ + [$\mu_0$] .5. Assume[$f_7$]($x_7$)= $K_7x_7$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_7(x_7)$= [0$x_7$]/(2 + $x_7$); the functional responses [[$f_7$]($x$) = [0] $x_7$/(1 + $x_7$)and $h_7(x_7)$=(1+$\omega_7x_7$)/(1 + $x_7$)(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_7 + [\mu_0] + [\mathrm{K}][P_0])}{[m_7]_7 + (\mu_0)_7 + [\mathrm{K}][P_0])_7^3\ [\mathrm{K}]_7[P_0]_7^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1+P_0) < b_7$ i e $[P_0] < b/(\delta_7 - b_7)$then the equilibrium$(P_0)$ $[P_0]_7(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_7$

$[f_{7*}] = K_7\ \mathrm{P}^*\ [f_{7*}] = \frac{\sigma \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*} \quad [h_*] = \frac{7+\omega_7 \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*}$

$[f_{7*}] = K_7\ [\mathrm{f}_{7*}^{\ 7}] = \frac{\sigma_7}{(7+\mathrm{Q}_7^*)^7} [\ \mathrm{g}_*^{\ 7}] = \frac{\sigma_7}{(7+\mathrm{P}_7^*)^7} \quad [\mathrm{h}_*^{\ 7}] = \frac{\omega_7 - 7}{(7+\mathrm{Q}_7^*)^7}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_7}{(m + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma)^7} + \frac{1\omega_7}{([m_7]_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma_7)^3} < 1$$

And

$$\frac{1}{(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^7} + \frac{1\omega_7}{(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $\mathrm{Q}^*$ which can be solved from the following equation:

$(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7} + (m_7 + [\mu_0]_7 + \frac{K_7 b_7}{\delta_7 - b_7} + \sigma_7)\ \mathrm{Q}_7^*)^3$

$= m_7 + [\mu_0]_7 + 3 + \frac{K_7 b_7}{\delta_7 - b_7}(1 + \mathrm{Q}^*)(1 + \omega_7 \mathrm{Q}_7^*$

Now

$$\left[\mathrm{a}\begin{matrix}1\\1\end{matrix}\right]-4\left[\mathrm{a}_{7}\right]\mathrm{e}=\left(\frac{1K_7 b_7}{\delta_7-b_7}\right)^{7}+\frac{4K_7 b_7\sigma_7\,\mathrm{Q}_{7}^{*}}{(\delta_7-b_7)(7+\mathrm{Q}_{7}^{*})}+\frac{\sigma_{7}^{7}\mathrm{Q}_{7}^{*7}}{\left(1+\mathrm{Q}_{7}^{*}\right)^{7}}$$

$$\frac{4\,b_7\sigma_7\,(\delta_{7}-\mathrm{b})\mathrm{Q}^{*}}{\delta_{7}\,(1+\mathrm{Q}_{7}^{*})^{7}}-\frac{4\,(\delta_{7}-b_7)^{7}\,(1-\omega_7)\mathrm{Q}_{7}^{*}}{\delta_{7}\left(1+\omega_7\mathrm{Q}_{7}^{*}\right)(1+\mathrm{Q}_{7}^{*})}\mathrm{P}*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_{7}\,ae_{7}^{mu}\;\;[\mu_0]_{7}-\frac{a_7}{a_7+7}$$

Hence if $[\mathrm{a}_{7}^{7}>4(a_7)_{7}\mathrm{e}$ and

In the following are given some more modular structures that are applied with the model and the concomitant results derived thereof:

Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments

(1) We have an intuitive understanding of how the world works. Our intuition works well in our everyday lives, but it doesn't quite match (e&eb) reality.

(2) One of the basic assumptions authors make about objects is that they are (=) objectively real.

(3) For example, when you leave your apartment in the morning, you imagine your house is right where you left it as you go about your day. Barring some calamity, you expect to find your apartment just as you left it when (e) you return in the evening.

(4) This idea is known as Macrorealism, in that a macroscopic object such as your house is still there even when (e) you're not looking at it

(5) In scientific terms we would say the state of an object exists independent of (e) an observer. While this seems obviously true, we can test the idea experimentally.

(6) To show this, let's consider an experiment. Suppose we have a box that contains a coin. At any time we can open the box and (eb) see whether it's heads or tails.

(7) The catch is (=) that the box has to follow a certain rule. Whatever you observe initially, (eb) you must observe the same thing ten seconds later.

(8) If you open the box and see heads, then (eb) ten seconds later when you open the box you must also see heads.

$$[a_7] = \frac{1K_7 b_7}{\delta_7 - b_7} + \frac{\sigma_7 \, \mathrm{Q}_7^*}{1 + \mathrm{Q}_7^*} > 1\vartriangle$$

This equation is important in its consequence in respect of following cases for which the model is applied:

**Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments**

(37) For the experiment we'll open the box, then open it five seconds later, then finally open it again at ten seconds. We know the first and last observations have to agree, but what about the observation in the middle? It depends upon what the box does with (e&eb) the coin.

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| (38) One possibility is that the box does nothing, so if you see tails the first time (eb), you will see tails both other times. |
| (39) Another possibility is that each time you open the box (eb) it flips the coin when you close it. |
| (40) So the first time you see heads, (eb) the second time tails, and (eb) the third time heads again. |
| (41) But they could also make a tricky box that flips (e&eb) the coin at seemingly random times, but does so in such a way that every ten seconds it repeats (eb) the pattern. In that case, authors would always see the same thing the first and last time, but the middle observation would seem (eb) random. |
| (42) If Macrorealism is true, then (eb) middle observations being always the same as (=) the first and last observation, always opposite or always random are (=) the only possibilities for the experiment |

**Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry Download as PDF (Size: 62KB) HTML PP. 50-53 DOI: 10.4236/jqis.2011.12007 5,508 Downloads 14,712 Views Citations Author(s) Leave a comment Mohamed S. El Naschie**

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| (37) Building upon the pioneering work of J. Bell [1] and an incredible result due to L. Hardy [2] it was shown (eb) that the probability of quantum entanglement of two particles is (=) a maximum of 9.0169945 percent [2]. |
| (38) This happens to be exactly the golden mean to (e) the power of five (?5) [3-7]. |
| (39) Although it has gone largely unnoticed for a long time, this result was essentially established independently in a much wider context by (e) the present author almost two decades ago [3-6]. |
| (40) The present work gives two fundamentally different derivations of (e) Hardy's beautiful result leading to (eb) precisely the same general conclusion, namely that by virtue of (e) the zero measure of the underlying |
| (41) Cantorian-fractal spacetime geometry the notion of (e) spatial separability in quantum physics is (=) devoid of any meaning [7]. The first derivation is purely logical and uses (e) a probability theory which combines the discrete with (e&eb) the continuum. |
| (42) The second derivation is purely geometrical and topological using (e) the fundamental equations of a theory developed by the author and his collaborators frequently referred to as E-infinity or Cantorian spacetime theory [3-7]. KEYWORDS Hardy's Quantum Entanglement, Golden mean, Cantor sets, Fractal Spacetime, E-Infinity Theory, Quantum Mechanics, J. S. Bell Cite this paper M. El Naschie, "Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry," Journal of Quantum Information Science, Vol. 1 No. 2, 2011, pp. 50-53. Doi: 10.4236/jqis.2011.12007. |

**Chaos, Solitons & Fractals Volume 28, Issue 4, May 2006, Pages 857–878 Sir Hermann Bondi 1919–2005 Cantorian spacetime and Hilbert space: Part I—Foundations G. Iovane http://dx.doi.org/10.1016/j.chaos.2005.08.074**

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| (61) Authors are going to show the link between the ε(∞) Cantorian space and( e&eb) the Hilbert spaces H(∞). |
| (62) In particular, El Naschie's ε (∞) is (=) a physical spacetime, i.e. an infinite dimensional fractal space, where (e) time is spacialized and the transfinite nature manifests (eb) itself. |
| (63) El Naschie's Cantorian spacetime is an arena where (e) the physics laws appear at each scale in a self-similar way linked to (e&eb) the resolution of the act of observation. |
| (64) By contrast the Hilbert space H (∞) is (=) a mathematical support, which describes (eb) the interaction between the observer and (e&eb) the dynamical system under (e&eb) measurement. |

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| (65) The present formulation, which is based on (e) the non-classical Cantorian geometry and topology of spacetime, automatically solves (e) the paradoxical outcome of (e) the two-slit experiment and the so-called particle–wave duality. |
| (66) In particular, measurement (i.e. the observation) is equivalent to (=) a projection of ε(∞) in the Hilbert space built on 3 + 1 Euclidean spacetime. |
| (67) Consequently, the wave–particle duality becomes (=) a mere natural consequence of (e) conducting an experiment in (eb) a spacetime with (e&eb) non-classical topological and (e&eb) geometrical structures, while (e) observing and taking measurements in (eb) a classical smooth 3 + 1 Euclidean spacetime. |
| (68) In other words, the experimental fact that a wave–particle duality exists is (=) an indirect confirmation of the existence of ε(∞) and a property of (e) the quantum–classical interface. |
| (69) Another direct consequence of the fact that real spacetime is) =) infinite dimensional hierarchical ε(∞) is (=) the existence of scaling law R(N), introduced by the author, which generalizes the Compton wavelength. |
| (70) It gives an answer to the problem of (e) segregation of matter at different scales, and shows (eb) the role of fundamental constants such as (=) the speed of light and Plank's constant h in the fundamental lengths scale without (e) invoking the principles of quantum mechanics. Copyright © 2005 Elsevier Ltd. All rights reserved |

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**Chaos, Solitons & Fractals Volume 14, Issue 9, December 2002, Pages 1341–1351 Fractal strings as an alternative justification for El Naschie's Cantorian spacetime and the fine structure constant Carlos Castro, http://dx.doi.org/10.1016/S0960-0779(02)00110-8**

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| (49) Beginning with the most general fractal strings/sprays construction recently expounded in the book by Lapidus and Frankenhuysen, it is shown how the complexified extension of El Naschie's cantorian-fractal spacetime model belongs to (e&eb) a very special class of families of fractal strings/sprays whose scaling ratios are given by (e) suitable pinary (pinary, p prime) powers of the Golden Mean. |
| (50) They then proceed to show why the logarithmic periodicity laws in (eb) Nature are direct physical consequences of (e) the complex dimensions associated with (e&eb) these fractal strings/sprays. |
| (51) Authors proceed with a discussion on quasi-crystals with (e&eb) p-adic internal symmetries, von Neumann's Continuous Geometry, the role of wild topology in (e&eb) fractal strings/sprays, (e&eb) the Banach-Tarski paradox, and tesselations of (e) the hyperbolic plane, quark confinement and (e&eb) the Mersenne-prime hierarchy of (e) bit-string physics in determining the fundamental physical constants in Nature. Fax: +1-914-965-3397 Copyright © 2002 Published by Elsevier Ltd. About ScienceDirectRemote access Shopping cart Contact and support Terms and conditions Privacy policy Cookies are used by this site. For more information, visit the cookies page. Copyright © 2016 Elsevier B.V. or its licensors or contributors. Science Direct ® is a registered trademark of Elsevier B.V. |

**Chaos, Solitons & Fractals Volume 21, Issue 2, July 2004, Pages 261–271 On Cantorian spacetime over number systems with division by zero Jakub Czajko, http://dx.doi.org/10.1016/j.chaos.2003.12.046**

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| (52) Division by zero is proposed in order to open (eb) new horizons for (e) abstract mathematics and physics. |
| (53) Although the formerly forbidden operation may be of (e) little practical value, the feasibility of unrestricted operations is needed for (e) meaningful deployment of linear vector spaces, whose duals are (=) inverse transposes of their primary vector spaces. |
| (54) Both quantum mechanics and (e&eb) pan-geometry of multispatial hyperspace operate on (e&eb) vectors that belong to (e&eb) mutually dual linear vector spaces. |
| (55) The division by zero also may remove (e) many years old paradox related to (e&eb) multiplication of zeros and makes it possible to introduce (eb) the concept of bigroup as (=) an additive and multiplicative group at the same time. |
| (56) This operational extension to (e&eb) algebraic groups makes thus (eb) an infinite-dimensional, fractal Cantorian spacetime (=) the preferred, invertible, operationally unrestricted, abstract mathematical infrastructure for physics. Tel.: +1-925-681-2058 Copyright © 2003 Elsevier Science Ltd. All rights reserved. |

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The nontrivial equilibrium is stable.

Some of the examples in which stability analysis is of seminal importance are: These are the system to which the

model becomes applicable

**Observation of a Dissipation-Induced Classical to Quantum Transition J. Raftery, D. Sadri, S. Schmidt, H. E. Türeci, and A. A. Houck Phys. Rev. X 4, 031043 – Published 8 September 2014**

(615)Here, authors report the experimental observation of (e) a dynamical quantum phase transition in (eb) a strongly interacting open photonic system.

(616) The system studied, comprising a Jaynes-Cummings dimer realized on (eb) a superconducting circuit platform, exhibits (eb) a dissipation-driven localization transition.

(617)Signatures of the transition in the homodyne signal and photon number reveal (eb) this transition to be from (e) a regime of classical oscillations into (e&eb) a macroscopically self-trapped state manifesting revivals, (=) a fundamentally quantum phenomenon.

(618) This experiment also demonstrates (eb) a small-scale realization of a new class of quantum simulator, whose well-controlled (e&eb) coherent and dissipative dynamics is suited to (e) the study of quantum many-body phenomena out of (e) equilibrium. Received 10 December 2013 DOI:http://dx.doi.org/10.1103/PhysRevX.4.031043 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**Viewpoint: Fermionic Vortices Find their Dual John McGreevy, Department of Physics, University of California at San Diego, La Jolla, CA 92093-0319, USA June 27, 2016• Physics 9, 72Theoretical work reveals a surprising relationship between the physics of fermionic vortices and quantum electrodynamics. Figure caption expand figure APS/Alan Stonebraker Figure 1: The action of discrete symmetries on the fermionic particle (top) and its dual fermionic vortex (bottom).**

(619)Dirac fermions in 2+1 dimensions (two spatial dimensions plus time) have been found in (eb) graphene and on (e&eb) the surfaces of topological insulators (TIs).

(620) A trio of papers now shows that these particles can also be found in (eb) a dual description of quantum electrodynamics (QED) [1–3].

(621)This deceptively simple statement joins (e&eb) together several disparate lines of progress in theoretical physics: interacting topological insulators, (e&eb) anomalies in quantum field theory, (e&eb) spin liquids, and (e&eb) compressible Quantum Hall States.

(622)Humans are not very good at understanding interacting many-body systems.

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| (623)We understand the behavior of free particles and can sometimes use (e) perturbation theory to (e) incorporate weak interactions between (e&eb) them. |
| (624)But what happens when we couple (e&eb) the particles strongly? |
| (625)To reveal the interesting collective phenomena that can result, we've had to rely largely on (e) experiments. |
| (626)Examples include (e) superconductivity (e&eb) superfluidity (e&eb) the integer and (e&eb) fractional quantum Hall effects (e&eb) the fractionalization of charge and spin, and (e&eb) the wetness of water. |
| (627)The fractional quantum Hall effect (FQHE) is a physical phenomenon in which the Hall conductance of 2D electrons shows (eb) precisely quantised plateaus at (eb) fractional values of (e) {\displaystyle e^{2}/h} e^{2}/h. |
| (628) It is a property of (e) a collective state in which electrons bind (e&eb) magnetic flux lines to make (eb) new quasiparticles, and excitations have (e) a fractional elementary charge and possibly also fractional statistics. The 1998 Nobel Prize in Physics was awarded to Robert Laughlin, Horst Störmer, and Daniel Tsui for the discovery and explanation of the fractional Hall effect.[1] |
| (629) However, Laughlin's explanation was (=) a phenomenological guess and only applies to (e&eb) fillings {\displaystyle \nu =1/m} \nu =1/m where (e) {\displaystyle m} m is an odd integer. The microscopic origin of the FQHE remains unknown and is a major research topic in condensed matter physics. (Wikipedia) |
| (630)The quantum Hall effect (or integer quantum Hall effect) is a quantum-mechanical version of (e) the Hall effect, observed in (eb) two-dimensional electron systems subjected to (e&eb) low temperatures and strong magnetic fields, in which the Hall conductance σ undergoes (e&eb) certain quantum Hall transitions to take on (e&eb) the quantized values {\displaystyle \sigma ={\frac {I_{\text{channel}}}{V_{\text{Hall}}}}=\nu \;{\frac {e^{2}}{h}},} \sigma = \frac{I_\text{channel}}{V_\text{Hall}} = \nu \; \frac{e^2}{h}, where {\displaystyle I_{\text{channel}}} I_\text{channel} is the channel current, {\displaystyle V_{\text{Hall}}} V_\text{Hall} is the Hall voltage, e is the elementary charge and h is Planck's constant. The prefactor ν is known as the "filling factor", and can take on either integer (ν = 1, 2, 3, ...) or fractional (ν = 1/3, 2/5, 3/7, 2/3, 3/5, 1/5, 2/9, 3/13, 5/2, 12/5, ...) values. |
| (631) The quantum Hall Effect is referred to as (=) the integer or fractional quantum Hall effect depending on whether ν is an integer or fraction, respectively. |
| (632)The striking feature of the integer quantum Hall Effect is (=) the persistence of the quantization (i.e. the Hall plateau) as (e) the electron density is varied. |
| (633) Since the latter remains constant when (e) the Fermi level is in a clean spectral gap, it must be that the Fermi level is immersed in (eb) essential spectrum (i.e. of finite density of states), though Anderson localized. |
| (634) The quantization and stability of the Hall conductance in such conditions was achieved in 1994 by (e) Bellissard, van Elst and Schulz-Baldes using the methods of (e) Non-Commutative Geometry.[1] |
| (635)The fractional quantum Hall effect is more complicated, as its existence relies fundamentally on (e) electron–electron interactions. |
| (636)Although the microscopic origins of the fractional quantum Hall effect are unknown, there are several phenomenological approaches that provide (eb) accurate approximations. |
| (637) For example the effect can be thought of as (=) an integer quantum Hall effect, not of electrons but of (e) charge-flux composites known as (=) composite fermions. |

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| (638) In 1988, it was proposed that there was quantum Hall effect without Landau levels. This quantum Hall Effect is referred to as (=) the quantum anomalous Hall (QAH) effect. |
| (639)There is also a new concept of the quantum spin Hall effect which is an analogue of the quantum Hall effect, where (e) spin currents flow instead of (e&eb) charge currents.[2] |
| (640)In two dimensions, when classical electrons are subjected to (e&eb) a magnetic field they follow (eb) circular cyclotron orbits. |
| (641)When the system is treated quantum mechanically (eb) these orbits are quantized. |
| (642)The energy levels of these quantized orbitals take on (e) discrete values |
| (643):{\displaystyle \textstyle E_{n}=\hbar \omega _{c}(n+1/2),} \textstyle E_{n}=\hbar \omega _{c}(n+1/2), where ωc = eB/m is the cyclotron frequency |
| (644) These orbitals are known as Landau levels, and at weak magnetic fields, their existence gives rise to (eb) many interesting "quantum oscillations" such as the Shubnikov–de Haas oscillations and the de Haas–van Alphen effect (which is often used to map the Fermi surface of metals). |
| (645) For strong magnetic fields, each Landau level is (=) highly degenerate (i.e. there are many single particle states which have the same energy En). |
| (646)Specifically, for a sample of area A, in magnetic field B, the degeneracy of each Landau level is (=) {\displaystyle \textstyle N=g_{s} BA/\phi _ {0},} \textstyle N=g_{s} BA/\phi _ {0}, where gs represents (eb) a factor of 2 for spin degeneracy, and ϕ0 = 2·10−15 Wb is the magnetic flux quantum. |
| (647) For sufficiently strong B-fields, each Landau level may have (e) so many states that all of the free electrons in the system sit in only a few Landau levels; it is in this regime where (e) one observes the quantum Hall Effect. |
| (648)The integers that appear in the Hall Effect are (=) examples of **topological quantum numbers**. |
| (649) They are known in mathematics as the first **Chern numbers** and are closely related to (e&eb) **Berry's phase.** |
| (650)A striking model of much interest in this context is the Azbel-Harper-Hofstadter model whose quantum phase diagram is (=) the Hofstadter butterfly shown in the figure. (deleted due to spatial constraints) |
| (651)The vertical axis is (=) the strength of the magnetic field and the horizontal axis is the chemical potential, which fixes the electron density. The colors represent the integer Hall conductance. Warm colors represent positive integers and cold colors negative integers. The phase diagram is fractal and has structure on all scales. In the figure there is an obvious self-similarity. |
| (652)Concerning physical mechanisms, impurities and/or particular states (e.g., edge currents) are important for (e) both the 'integer' and 'fractional' effects. |
| (653) In addition, Coulomb interaction is also essential in (eb) the fractional quantum Hall effect. |
| (654)The observed strong similarity between integer and (=) fractional quantum Hall effects is explained by (e) the tendency of electrons to form bound states with (e&eb) an even number of magnetic flux quanta, called composite fermions. |
| (655)Duality is the circumstance in which cranking up interactions results in (eb) another system we understand, usually one that can be described in terms of (e&eb) another set of weakly coupled particles. |
| (656)When such a dual representation is found, (eb) it gives a solution to the strong coupling problem. |
| (657) A further reason that theoretical physicists get excited about dualities is that they undermine (e) the very notion of an “elementary particle.” |
| (658)In terms of the original variables, the dual, weakly coupled particles are (=) generally solitons—large, heavy, classical objects made up of (e) many quanta. |

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| (659)The remarkable thing is that solitons can become (=) "elementary," in the sense that the model that describes the system can be reformulated with the solitons as (=) the basic constituents. |
| (660)An example of a duality, which we will use below, arises in (eb) Maxwell's theory of electromagnetism, where (e) the strength of the interactions is determined by (e) the electric charge e. |
| (661)The duality exchanges electric and (e&eb) magnetic point charges, and the interaction strength in the dual description is 1⁄e, which is small when (e) e is large. |
| (662)This idea has been generalized to (e) many other field theories, thanks, in particular, to tools of (e) supersymmetry [4]. |
| (663)Another classic example of duality involves (e&eb) vortices. |
| (664) A vortex is a localized object in the plane around which a phase field winds (e&eb). (A phase field takes values on the circle.) |
| (665)Dualities relating particles and vortices have (e) a rich history in particle physics and condensed-matter physics. |
| (666) The simplest version relates the particles in the normal state of a superfluid to (e&eb) the vortices in a dual superconductor [5]. |
| (667)This charge-vortex duality has been (=) a workhorse in the study of two-dimensional condensed-matter systems, such as (=) fractional Quantum Hall States, superconducting thin films, and beyond-Landau critical phenomena [6]. |
| (668)Vortices can be (=) fermions (as they frequently are, for instance, in type-II superconductors). |
| (669) A dual description of a system with (e&eb) fermionic vortices thus requires (e) a new ingredient, since (e) it must include (e) such a fermionic particle. |
| (670)The new work [1–3] builds on (e&eb) earlier theoretical work [7] to make (eb) an appealing proposal for (e) the required duality: the fermionic vortex of QED in (eb) 2+1 dimensions is (=) a free Dirac fermion. The argument for this proposal builds on our understanding of TIs. |
| (671)TIs are distinguished only by (e) their edge states. |
| (672) The most famous example is a (3+1)-dimensional time-reversal invariant TI, which can host boundary states whose band structure is defined by (e) a single Dirac cone. |
| (673)A single time-reversal-symmetric Dirac cone is not possible without (e) the bulk TI, so its appearance in experiments provides (eb) a signature of the bulk phase. |
| (674) But the Dirac cone is not (e) the only possible signature. |
| (675) Deformations of the boundary conditions can dramatically change (e&eb) the spectrum of the edge states, yet preserve some essential "TI-ness." This essential property is (=) called an anomaly in high-energy theory. |
| (676)Max Metlitski of the University of California, Santa Barbara, and Ashvin Vishwanath of the University of California, Berkeley [1], and independently, Chong Wang and Todadri Senthil at the Massachusetts Institute of Technology, Cambridge [2], use this ambiguity in the edge theory and the electric-magnetic duality of an auxiliary theory of bulk electromagnetism coupled to (e&eb) the TI to argue for the duality. |
| (677)When I said that (3+1)-dimensional QED had a duality that interchanged the electric and magnetic charges, one could have (e) complained that the electron is (=) a fermion, while (e) a magnetic monopole may not be. |
| (678)However, the authors observe that in a TI, a charge-two monopole is (=) also a fermion. |

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| (679) By analyzing the spectrum of electric and magnetic charges, they conclude that the duality exchanges the electron and (e&eb) the charge-two monopole and, moreover, exchanges a TI with (e&eb) time-reversal symmetry with (e&eb) a "chiral TI," which instead has (e) particle-hole symmetry |
| (680). But this is just what is required of (e) a symmetry that exchanges a fermionic particle with (e&eb) a fermionic vortex: time-reversal symmetry maps the electron to itself but maps the vortex to (e&eb) another vortex whose phase winds (e&eb) the other way (see Fig. 1). |
| (681)The exchange of particles and holes, in turn, preserves (eb) the vortex. |
| (682)The bulk duality implies (eb) a duality between edge theories. |
| (683)The freedom to deform the edge of the TI implies (eb) a weak form of equivalence: although the edge theories are (=) topologically the same (in terms of symmetry assignments and (e&eb) anomalies), they need not have (e) identical low-energy physics. |
| (684)One side of the duality is (=) a free, massless Dirac particle, but the long-distance physics of (2+1)-dimensional QED is (=) a long-standing problem, which is important in the study of spin liquids. |
| (685) Many gauge theories in two dimensions develop (eb) a gap in their energy spectrum because of (e) the proliferation of monopoles. |
| (686) In the realization proposed by Metlitski and Vishwanath [1] and by Wang and Senthil [2], this cannot happen because (e) the monopoles carry electric charge, and indeed the stronger notion of duality as (=) an equivalence of low-energy physics may hold. |
| (687)Evidence for the stronger relation comes from (e) David Mross and colleagues at the California Institute of Technology, Pasadena [3], who have derived a lattice version of the duality using (e) a clever but explicit change of variables. |
| (688)Their construction is particularly interesting because (e) it uses (e) a lattice of one-dimensional models which themselves can only occur at the edge of a TI. |
| (689)Authors explicitly derive the duality between a free electronic Dirac cone and (e&eb) quantum electrodynamics in (eb) (2+1) dimensions (QED3) with (e&eb) N=1 fermion flavors. |
| (690)The duality proceeds via (e&eb) an exact, nonlocal mapping from electrons to (e&eb) dual fermions with (e&eb) long-range interactions encoded by (e) an emergent gauge field. |
| (691)This mapping allows (eb) us to construct (eb) parent Hamiltonians for exotic topological-insulator surface phases, derive (eb) the particle-hole-symmetric field theory of (e) a half-filled Landau level, and nontrivially constrain (e) QED3 scaling dimensions. |
| (692)They similarly establish duality between bosonic topological insulator surfaces and (e&eb) N=2 QED3. |

**Black hole thermodynamics with generalized uncertainty principle Li Xiang, X. Q. Wen (Submitted on 6 Jan 2009 (v1), last revised 18 Mar 2009 (this version, v2))**

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| (693) In the standard viewpoint, the temperature of a stationary black hole is proportional to (=) its surface gravity, $T_H=\hbar\kappa/2\pi$. |
| (694)This is a semiclassical result and the quantum gravity effects are not (e) taken into consideration. |
| (695)This Letter explores a unified expression for (e) the black hole temperature in the sense of (e) a generalized uncertainty principle (GUP). |
| (696) Discussion involves a heuristic analysis of (e) a particle which is absorbed by (e) the black hole. |
| (697)Besides a class of static and spherically symmetric black holes, an axially symmetric Kerr-Newman black hole is considered. Different from the existing literature, authors suggest (eb) that the black hole's irreducible mass represent (eb) the characteristic size in (eb) the absorption process. |

(698)The information capacity of a remnant is also discussed by (e) Bousso's D-bound in de Sitter spacetime. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP 0910:046,2009 DOI: 10.1088/1126-6708/2009/10/046 Cite as: arXiv:0901.0603 [gr-qc] (or arXiv:0901.0603v2 [gr-qc] for this version)

IF we define

$$[a_7]_7 = [f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}]\,P_7^*$$

$$[a_7]_7 = \left[f_{7*}^{7}\right] P_7^* - \frac{h_*^7}{h_*} P^*(0)$$

$$e = [g_*^7]_7\, Q_7^*$$

$$\eta_7(t) = \left[h_*\right]_7 \beta_7 \left(\theta_7(t)\right) - \mu_7(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\left[a_7\right] M_7 - [a_7]_7 W_7 - [n(t)] M_7$$

$$w_7(t) = eM_7$$

Write

$$A = \begin{pmatrix} -[a_7]_7 - [a_7]_7 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B = (t) = \begin{pmatrix} [\eta(t)]_7 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_7 \\ W_7 \end{pmatrix}$$

The solution of (32)is

$$X_7(t) = e^{A,t}[X]_7(0) + \int_0 e^{At}$$

## PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION

## ATTRITION MODELS

## MODULE EIGHT

Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal.

### 61. MODEL FORMULATION

**Model is applied to systems given in the following:**

**Researchers blur the line between classical and quantum physics by connecting chaos and entanglement July 12, 2016 by Sonia Fernandez**

(3641) Experimental link between quantum entanglement (left) and (e&eb) classical chaos (right) found using (e) a small quantum computer. Credit: University of California - Santa Barbara Using a small

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| quantum system consisting of three superconducting qubits, researchers at UC Santa Barbara and Google have uncovered a link between aspects of classical and (e&eb) quantum physics thought to be unrelated: classical chaos and (e&eb) quantum entanglement. |
| (3642) Their findings suggest (eb) that it would be possible to use (e) controllable quantum systems to investigate certain fundamental aspects of nature. |
| (3643) "It's kind of surprising because (e) chaos is this totally classical concept—there's (=) no idea of chaos in a quantum system," Charles Neill, a researcher in the UCSB Department of Physics and lead author of a paper that appears in Nature Physics. |
| (3644) "Similarly, there's no concept of entanglement within (eb) classical systems. And yet it turns out that chaos and (e&eb) entanglement are really very strongly and clearly related." |
| (3645) Initiated in the 15th century, classical physics generally examines and describes (eb) systems larger than atoms and molecules. |
| (3646) It consists of hundreds of years' worth of study including Newton's laws of motion, electrodynamics, relativity, thermodynamics as well as chaos theory—the field that studies the behavior of (e) highly sensitive and unpredictable systems. |
| (3647) One classic example of chaos theory is the weather, in which a relatively small change in one part of (e) the system is enough to foil predictions—and vacation plans—anywhere on the globe. |
| (3648) At smaller size and length scales in nature, however, such as those involving (e&eb) atoms and photons and their behaviors, classical physics falls short. |
| (3649) In the early 20th century quantum physics emerged, with its seemingly counterintuitive and sometimes controversial science, including (e) the notions of superposition (the theory that a particle can be located in (eb) several places at once) and entanglement (particles that are deeply linked behave as (=) such despite physical distance from one another). |
| (3650) And so began the continuing search for connections between the two fields. All systems are fundamentally quantum systems, according Neill, but the means of describing in a quantum sense the chaotic behavior of, (e) say, air molecules in an evacuated room, remains limited. |
| (3651) Imagine taking a balloon full of air molecules, somehow tagging them so you could see them and then releasing them (eb) into a room with no air molecules, noted co-author and UCSB/Google researcher Pedram Roushan. |
| (3652) One possible outcome is that the air molecules remain clumped together (e&eb) in a little cloud following the same trajectory around the room |
| (3653) And yet, he continued, as we can probably intuit, the molecules will more likely take off in (eb) a variety of velocities and directions, bouncing off walls and interacting with (e&eb) each other, resting after the room is (=) sufficiently saturated with them. |
| (3654) "The underlying physics is (=) chaos, essentially," he said. |
| (3655) The molecules coming to rest—at least on (eb) the macroscopic level—is the result of (e) thermalization, or of reaching equilibrium after they have achieved (eb) uniform saturation within the system. |
| (3656) But in the infinitesimal world of quantum physics, there is (=) still little to describe that behavior. |
| (3657) The mathematics of quantum mechanics, Roushan said, do not allow for (e) the chaos described by Newtonian laws of motion. |
| (3658) To investigate, the researchers devised an experiment using (e) three quantum bits, the basic computational units of (e) the quantum computer. |

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| (3659) Unlike classical computer bits, which utilize (e) a binary system of two possible states (e.g., zero/one), a qubit can also use (e) a superposition of both states (zero and one) as a single state. |
| (3660) Additionally, multiple qubits can entangle, or link so closely that their measurements will (=) automatically correlate. |
| (3661) By manipulating these qubits with electronic pulses, Neill caused them to interact, rotate and evolve (e&eb, e&eb, eb) in the quantum analog of (e) a highly sensitive classical system. |
| (3662) The result is a map of entanglement entropy of (e) a qubit that, over time, comes to strongly resemble that of (e) classical dynamics—the regions of entanglement in (eb) the quantum map resemble the regions of (e) chaos on the classical map. |
| (3663) The islands of low entanglement in (eb) the quantum map are located in the places of (e) low chaos on the classical map. |
| (3664) "There's a very clear connection between entanglement and (e&eb) chaos in these two pictures," said Neill. |
| (3665) "And, it turns out that thermalization is (=) the thing that connects chaos and (e&eb) entanglement. |
| (3666) It turns out that they are actually the driving forces behind (eb) thermalization. |
| (3667) "What we realize is that in almost any quantum system, including on quantum computers, if you just let it evolve and you start to study what happens as a function of time, (eb) it's going to thermalize," added Neill, referring to the quantum-level equilibration. "And this really ties together the intuition between classical thermalization and (e&eb) chaos and how it occurs (eb) in quantum systems that entangle." |
| (3668) The study's findings have fundamental implications for (e) quantum computing. |
| (3669) At the level of three qubits, the computation is relatively simple, said Roushan, but as researchers push to build increasingly sophisticated and powerful quantum computers that incorporate (e) more qubits to study highly complex problems that are beyond the ability of classical computing—such as those in the realms of machine learning, artificial intelligence, fluid dynamics or chemistry—a quantum processor optimized for such calculations will be a very powerful tool. |
| (3670) "It means we can study things that are completely impossible to study right now, once we get to bigger systems," said Neill.More information: Nature Physics, DOI: 10.1038/nphys3830 , http://www.nature.com/nphys/journal/vaop/ncurrent/full/nphys3830.html Journal reference: Nature Physics search and more info website Provided by: University of California - Santa Barbara |

**Title: coherent control of macroscopic quantum states in a single-Cooper-pair box Authors: Nakamura, Y.; Pashkin, Yu. A.; Tsai, J. S. Publication: Nature, Volume 398, Issue 6730, pp. 786-788 (1999). (Nature Homepage) Publication Date: 04/1999 Origin: NATURE DOI: 10.1038/19718 Bibliographic Code: 1999Natur.398...786N**

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| (3671) A nanometre-scale superconducting electrode connected to (e&eb) a reservoir via (e&eb) a Josephson junction constitutes (e) an artificial two-level electronic system: a single-Cooper-pair box. |
| (3672) The two levels consist of (e) charge states (differing by 2e, where (e) e is the electronic charge) that are coupled by (e) tunnelling of Cooper pairs through (e&eb) the junction. |
| (3673) Although the two-level system is macroscopic, containing (e) a large number of electrons, the two charge states can be (=) coherently superposed. |

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| (3674) | The Cooper-pair box has therefore been suggested as (=) a candidate for a quantum bit or `qubit'- the basic component of a quantum computer. |
| (3675) | Here authors report the observation of quantum oscillations in (eb) a single-Cooper-pair box. |
| (3676) | By applying a short voltage pulse via (e&eb) a gate electrode, we can control (e&eb) the coherent quantum state evolution: the pulse modifies (e&eb) the energies of (e) the two charge states non-adiabatically, bringing them into (e&eb) resonance. |
| (3677) | The resulting state-a superposition of (e) the two charge states-is detected by (e) a tunnelling current through (e&eb) a probe junction. |
| (3678) | Results demonstrate electrical coherent control of (e) a qubit in a solid-state electronic device. |

**Superconducting quantum computing From Wikipedia, the free encyclopedia**

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| (3679) | Superconducting quantum computing is a promising implementation of (e) quantum information technology that involves (e&eb) nanofabricated superconducting electrodes coupled through (e&eb) Josephson junctions. |
| (3680) | As in a superconducting electrode, the phase and (e&eb) the charge are conjugate variables. |
| (3681) | There exist three families of superconducting qubits, depending on (=) whether the charge, (e&eb) the phase, or neither of (e) the two is good quantum number. |
| (3682) | These are respectively termed charge qubits (e&eb) flux qubits, and (e&eb) hybrid qubits. |

**Theory**

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| (3683) | Unlike many other physical implementations of a qubit which involve (e&eb) exclusively two-level systems (such as nuclear spin and (e&eb) photon polarization), the integrated quantum circuit involved in (e&eb) a superconducting qubit is (=) a multi-level system, of which only the first two levels are used as (=) the computational basis. |
| (3684) | A basic requirement for such an implementation is that the energy levels are (=) not uniformly spaced, so that (e) photons of a particular frequency which cause (eb) transition between the 0 and 1 levels do not(e) cause transitions from the first level to (e&eb) the higher levels as well. |
| (3685) | For electronic signals to be carried from one part of the circuit to (e&eb) another without energy loss and (e&eb) hence decoherence, the system also needs to be (=) non-dissipative, i.e., the metallic parts involved should have (e) zero resistance. |
| (3686) | These circuits currently need to be (=) operated at very low temperatures so that, 1) superconductivity is (=) realized, and 2) thermal fluctuations do not (e) cause transitions between (e&eb) energy levels. |

**Quantum LC Circuit**

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| (3687) | The simplest non-dissipative quantum circuit consists simply of (e) an inductor L {\displaystyle L} L and (e&eb) capacitor C {\displaystyle C} C, with (e&eb) the metallic wires connecting them being superconducting. |
| (3688) | Here the flux Φ {\displaystyle \Phi } \Phi through (e&eb) the inductor and (e&eb) the charge Q {\displaystyle Q} Q in the capacitor are (=) canonically conjugate variables which obey (eb) the commutation relation [ Φ , Q ] = i ℏ {\displaystyle [\Phi ,Q]=i\hbar } [\Phi ,Q]=i\hbar |

**The Hamiltonian associated with this system is**

(3689) H = Φ 2 2 L + Q 2 2 C {\displaystyle H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}} H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}

(3690) This gives rise to the energy levels E = ℏ ω 0 ( n + 1 2 ) {\displaystyle E=\hbar \omega _{0}(n+{\frac {1}{2}})} E=\hbar \omega _{0}(n+{\frac {1}{2}}) Here ω 0 = 1 L C {\displaystyle \omega _{0}={\frac {1}{\sqrt {LC}}}} \omega _{0}={\frac {1}{{\sqrt {LC}}}}

**Quantum gravity Carlo Rovelli (2008), Scholarpedia, 3(5):7117. doi:10.4249/scholarpedia.7117 revision #135539 [link to/cite this article] Post-publication activity Curator: Carlo Rovelli Dr. Carlo Rovelli, Center of Theoretical Physics, and Marseille, France**

(3691) Quantum Gravity is the name given to any theory that describes gravity in (eb) the regimes where quantum effects cannot be disregarded.

(3692) At present, there is no such a theory which is universally accepted and confirmed by experience. Therefore the term "Quantum Gravity" indicates (eb) more an open problem than a specific theory.

(3693) Several research lines, at different levels of development, offer tentative solutions to the problem. These tentative quantum-gravity theories are variously viewed as competing research directions, or as contributions to the common goal of finding the physically correct theory. The quest for the good quantum theory of gravity bears on a number of fundamental issues and it is sometimes presented as the most important open problem in fundamental physics: the "Holy Grail" of contemporary theoretical physics.

**The problem of quantum gravity**

(3694) The current theory of (e) gravity is general relativity, a theory developed by Albert Einstein at the beginning of the XX century, which has changed in depth our understanding of (e) what are space and time.

(3695) Similarly, quantum mechanics has changed in depth our understanding of matter, energy and causality. These two theories, today very well empirically supported, have opened (eb) a major conceptual revolution in physics.

(3696) However, they appear to be incompatible, at least at a first reading, because (e) each of the two is formulated on the basis of (e) principles that are explicitly contradicted by (e) the other theory.

(3697) Therefore in spite of the immense development of scientific knowledge it has delivered, the physics of the XX century has left us with (e&eb) a major confusion on what are the basic conceptual ingredients for (e) understanding the physical world.

(3698) In other words, the scientific revolution opened by general relativity and quantum mechanics at the beginning of (e) the XX century is not concluded yet, and a new synthesis is required. Quantum gravity, merging general relativity and quantum mechanics should be this new synthesis.

**Quantum space and quantum time**

(3699) Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.

(3700) In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.

(3701) They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.

(3702) When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.

(3703) Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.

**The Planck scale**

(3704) Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".

(3705) Here $\hbar$ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.

(3706) The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.

(3707) Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)

(3708) Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.

## Physical applications

Some phenomena where quantum gravity plays a major role are the following.

**The microscopic structure of spacetime**

(3709) If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.

(3710) Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length

(3711) Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.

(3712) A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.

**Early cosmology**

(3713) According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.

(3714) However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.

(3715) A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.

(3716) Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".

(3717) One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.

**Black holes**

(3718) Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.

(3719) First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.

(3720) Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.

(3721) What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.

(3722) Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.

**Astrophysical effects**

(3723) Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.

(3724) An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.

(3725) The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.

**Main approaches**

(3726) The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.

(3727) Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.

(3728) The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.

**String theory**

(3729) String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.

(3730) Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.

(3731) String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.

(3732) Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.

(3733) However, when summed the entire series appears to be (=) divergent.

(3734) In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.

(3735) Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.

(3736) Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.

(3737) In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.

**Loop quantum gravity and spin foams**

(3738) Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.

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| (3739) The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity. |
| (3740) Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime. |
| (3741) The granular properties of space can be (=) explicitly computed. |
| (3742) In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory. |
| (3743) The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks. |
| (3744) Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure. |
| (3745) The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit. |
| (3746) Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism. |
| (3747) This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism. |

**Noncommutative geometry**

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| (3748) Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime. |
| (3749) Quantum mechanics teaches us that dynamical quantities are (=) noncommutative --meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed. |
| (3750) It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry. |
| (3751) A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces. |
| (3752) Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces. |
| (3753) This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity. |

**Others**

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| (3754) Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays |

(3755) . Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection.

(3756) Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible.

(3757) The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.

Let $p_8(t, a_8)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(g_8)Q_8$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_8(0) = 0, \quad g_8(P_8) > 0, \quad P_8 \text{ 0}, \qquad \lim_{P_8 \to \infty} g_8(P_8) = (K_8)_8 \text{ }\infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_8 + (\rho_a)_8 + D_8(a_8, P_8, Q_8)\rho_8 + 0, \qquad 0 < a_8 \text{ }\infty,$$

$$\rho_8(t,0) = \int_0^\infty n_8(a_8, P_8, Q_8)\rho_8(t_8, a_8)da_8$$

Where $D_8$ is the mortality function and $n_8$ is the fertility function.

Put $_8$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_8}{dt}\right]_8 = x_8\left[\frac{dQ_8}{dt}\right]_8 - x_8\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_8}{dt}\right]_8 = x_3\left[\frac{dQ_3}{dt}\right]_8 - x_4\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_8}{dt}\right]_8 - x_6\left[\frac{dQ_3}{dt}\right]_8$$

Where $x_8, x_8, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_8$ and $g_8(P_8)Q_8$

Now when you add the term $(b_8Q_8)$ with $g_8(P_8)(Q_8)$ we are infect see is that $(b_8Q_8)$ is credited incrementally part of $g_8(P_8)(Q_8)$ or instant nearly, when subtracted from $\frac{dQ_8}{dt}$ , it means that $g_8(P_8)(Q_8)$ is slowly eaten by $\frac{dQ_8}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_8 = \mu_8(a_8) + (f_8)_8(P_8)_8 + (f_8)_8(Q)_8 \quad (2)$$

$$n_8 = \beta_8(a_8)h_8(Q_8), \quad (3)$$

Where $\mu_8$the age- dependent debit rate whit is $\mu_8(a_8)$ 0 $\forall a > 0$; $(f_8)_8$ is the density- dependent death rate with $f_8(0) = 0$, $(f_8^8)_8(P_8)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_8)_8(P_8) = (k_8)_8$ ; $f_8$ is the functional response of creditors to prey with $[f_8]_8(0) = 0, (f_8^8)_8(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_8)_8(Q_8) = [k_8]$ ; $\beta$ is the age

dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^8(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_8 + (\rho_a)_8 = -\mu_8(a_8) + (f_8)_8 P_8) + \big(f_8)_8(Q_8)\big)(\rho_8)_8 \quad (4)$$

$$\rho_8(t_8, 0) = h_8(Q_8)\int_0^\infty \beta_8(a_8)\rho_8(t, a_8)da, \quad (5)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(P_8)Q_8 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_8$ and $Q_8$ of (4-6) satisfy

$$0 \ \rho_8(t, a_8) \ \rho_8^*(t, a_8),$$

$$0 \ Q_8(t) \ Q_8^*(t),$$

Where $\rho_8^*(t, a_8)$ and $Q_8^*(t)$, are solutions of

$$(\rho_t^*)_8 + (\rho_a^*)_8 = -(\mu_8(a_8) + (f_8)_8(P_8^*) + (f_8)_8 Q_8^*))\,\rho_8^*, \quad (7)$$

$$\rho_8^*(t, 0) = \int_0^\infty \beta_8(a_8)\,\rho_8^*(t, a_8)da_8 \quad (8)$$

$$\frac{dQ_8^*}{dt} = -b_8 Q_8^* + g_8(P_8^*)Q_8^*, \quad (9)$$

With $P_8^*(t) = \int_0^\infty \rho_8^*\ (t, a_8)da_8$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_8^*(t)}{dt} = P_8^*(t)\big(r_8 - (f_8)(P_8^*) - (f_8)_8(Q^*)\big), \quad (10)$$

$$\frac{dQ_8^*(t)}{dt} = Q_8^*(t)(-b_8 + g_8(P_8^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_8(a)e^{-\int_0^a \mu_8(\tau)dt}e^{-ra}\ da_8 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_8^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(t)dt}\ da_8 > 1, \quad (H1)$$

Which guarantees $r_8 > 0$.

*Remark.* If $\beta_8(a) = [\beta_0]_8 a_8 e^{-\alpha_8 a_8}, \mu_8(a_8) \equiv \mu_8$ then the solution of (12) is

$$r_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_8]_8(P_8) = 0, [f_8](Q_8) = \lambda_8 Q_8\ h_8(Q_8) = 1, \quad and\ \ g_8(P_8) = c_8 P_8$$

$$[\rho_8]_8 + [\rho_a]_8 = -(\mu_8 + \lambda_8 Q_8)\rho_8$$

$$\rho_8(t,0) = \int_0^{\infty} [\beta_0]_8 a_8 e^{-\alpha_8 a_8} \rho_8(t,a_8)\, da_8$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + c_8 P_8 Q_8$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_8}{dt} = \omega_8 P_8 - \lambda_8 P_8 Q_8$$

$$\frac{dQ_8}{dt} = b_8 Q_8 - c_8 P_8 Q_8$$

Where $\omega_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$

## 62. THE EQUILIBRIUM $(\rho_0)_8(a_8), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_8(a_8), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_8 + (\rho_a)_8 = -\big(\mu_8(a_8) + (f_8)_8(P_8)\big)\rho,$$

$$\rho_8(t,0) = \int_0^{\infty} \beta_8(a_8)\rho_8(t,a_8)da_8 \qquad (13)$$

The solution of (13) has the form of

$$\rho_8(t,a_8) = Ce_8^{r(t-a_8)} - \int_0^a (\tau)d\tau - \int_0^8 (f_8)_8\big(P_8(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_8(t) = C_8 e^{r_8 t - \int_0^t (f_8)_8\left(P_8(\tau)\right)d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau}\, da.$$

Hence, there exists an equilibrium $P_0 = \int_8^{-8}(r)$ if $k_8 > r$. This is equivalent to

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{(k_8)_8 a_8} da_8 < 1.$$

In fact, if we define

$$R_8(s_8) = \int_0^\infty \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau)d\tau} e^{-s_8 a_8}\, da_8$$

$R(s)$ is a decreasing function. $R_8(r_8) = 1$ And $R_8(k_8)_8 < 1$ imply $r_8 < (k_8)_8$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_8(a_8), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{k_8 a_8} da_8 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_8 \equiv \int_0^\infty (\rho_0)_{88}(a_8)da_8$ *satisfies*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{-(f_8)_8(P_0)a_8} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_8(a_8), 0)$

Let $u_8(t,a_8) = \rho(t,a_8) - \rho_0(a_8)$ Then

$$N_8(t) \equiv \int_0^\infty u_8(t,a_8)da_8 = \int_0^\infty \rho_8(t,a_8)da_8 - \int_0^\infty (\rho_0)_8(a_8)da = P_8(t) - P_0)_8$$

And

$$(u_8)_8 + (u_a)_8 = (\rho_t)_8 + (\rho_a)_8 - (\rho_0^8)_8(a_8)$$

$$= -\mu_8(a_8) + (f_8)_8(P_8) + (f_8)_8(Q_8))\,(u_8 + (\rho_0)_8) + \mu_8(a_8)(\rho_0)_8 + (f_8)_8(P_0)_8(\rho_0)_8$$

$$= -\mu_8(a_8)u_8 - (f_8)_8(P_8) + (f_8)_8(Q_8))\,u_8 - (\rho_0)_8\,(f_8)_8(P_8) - (f_8)_8(P_0)_8$$

$$+ (f_8)_8(Q_8)), \qquad (15)$$

Where $(\rho_0^8)_8(a_8)$ denotes $d(\rho_0)_8(a_8)/da_8$

Expansion $(f_8)_8(P_8)$ and $(f_8)_8(Q_8)$ at $(P_0)_8$ and 0 yield respectively

$$(f_8)_8(P_8) = (f_8)_8((P_0)_8 + N_8)$$

$$= (f_8)_8(P_0)_8 + (f_8^8)_8(P_0)_8N_8 + o_8(N_8), \qquad (16)$$

$$(f_8)_8(Q_8)$$

$$= (f_8^8)_8(0)Q + o_8(Q_8) \qquad (17)$$

Write $f_8(P_0) = f_{80}, f_8'(P_0) = f_{80}'$, and $f_8'(0) = f_{80}'$ and substitute (16) and (17) into (15) to obtain

$$(u_8)_8 + (u_a)_8 \approx -(\mu_8(a_8) + (f_{80})_8)u - (f_{80}^8)_8(\rho_0)_8(N_8) - (f_{80}^8)_8(\rho_0)_8(Q_8)$$

$$u_8(t,0) = \rho_8(t,0)(\rho_0)_8(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a) + (\rho_0)_8(a_8)da_8 - (\rho_0)_8\,(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a_8)da_8 + (h_8(Q_8) - 1)(\rho_0)_8\,(0) \qquad (18)$$

Expanding

$$h_8(Q_8) = 1 + (h^8)_8(0)Q_8 + 0(Q_8) \equiv 1 + h_0^{'}Q_8 + 0Q_8$$

And substituting it into (18), we arrive at

$$u_8(t,0) \approx \int_0^\infty \beta_8\,(a_8)u_8(t,a_8)da_8 + (h_{a_8}^{'})(\rho_0)_8 Q_8$$

The linearization is the following :

$$[u_8] + [u_8] = -(\mu_8(a_8) + [t_{80}]_8 u_8 - [\alpha_8]_8(a_8)N_8 - [\alpha_8]_8(a_8)Q_8$$

$$u_8(t,0) = \int_0^\infty \beta_8(a_8)u_8(t,a_8)da_8 + [\xi_8]Q_8 \qquad (19)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8$$

where

$$[\alpha_8]_8(a_8) = [t_8^8 a]_8[p,o]_8(a_8)d2(a_8) = [f_{80}^8]_8(\rho\, o\,)_8(a_8)\xi = [\mathrm{h\ o}^8]_8[\rho \mathrm{o}]_8(\mathrm{o}) \text{ and } \mathrm{b}_8 = \mathrm{b}_8 - g_8(P_8)$$

If $g_8(P_0) > b$, clearly, $(\rho_0)_8(a_8), 0)$ is unstable and the predator population is persistent. Assume $g_8(P_0)_8 < b$.

Then, $Q_8(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_8 + u_a = -(\mu_8(a_8) + (f_{80})_8)u_8 - (\alpha_8)_8(a_8)N_8$$

$$= -(\mu_8(a_8) + f_{80})u_8 - \alpha_8(a_8)\int_0^\infty u_8(t,a_8)da_8$$

$$u_8(t,0) = \int_0^\infty \beta_8\,(a_8)u_8(t,a_8)da_8$$

Using the method developed by Cushing [2, 3], we suppose

$$u_8(t,a_8) = \tilde{u}_8(a_8)^{e^{\gamma(t-a_8)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_8}\,da_8$. Then,

A solution of the first differential equation for $\tilde{u}(a_8)$ gives,

$$\tilde{u}_8(a_8) = e^{-\int_0^a \mu_8(\tau_8)d\tau - (f_{80})_8 a}\left(\tilde{u}_8(0) - w\int_0^a e^{\int_0^\tau \mu_8(s_8)ds + (f_{80})}\tau_8\,(\alpha_8)_8(\tau_8)e^{\gamma_8\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\right)\tilde{u}_8(0)$$

$$+w\int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\int_0^a (\alpha_8)_8(\tau_8)e^{-\gamma_8\tau_8}e^{(f_{80})_8\tau_8 + \int_0^\tau \mu_8(s_8)ds}d\tau_8\,da_8$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}da\right)(\tilde{u}_8(0))_8$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_8 a_8-(f_{80})_8 a_8-\int_0^a \mu_8(\tau_8)d\tau}\int_0^a e^{(f_{80})_8\tau_8+\int_0^\tau \mu_8(s_8)ds+y\tau}\alpha_8(\tau)d\tau\, da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_8 a_8-(f_{80})_8 a_8-\int_0^a \mu_8(t_8)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_8\,(a_8)Y_8(\gamma_8,a_8)da_8\right)u_8(0)$$

$$+w\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)\int_0^a(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8=0,$$

(24)

$$\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\tilde{u}_8(0)$$

$$+\left(-1-\int_0^{\infty}Y_8(\gamma_8,a_8)\int_0^a(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8\right)w=0.$$

Equation (20) has a solution of the form of $u_8(t,a_8)=u_8(a_8)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_8(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_8(\gamma_8,a_8)=\int_0^a(\rho_0)_8(\tau_8)Y^{-8}\,(\gamma_8,a_8)d\tau_8$. We obtain the characteristic equation

$$f_{80_8}^{l}\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\left(\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$+\left(1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)da\right)\left(1+(f_{80}^{l})_8\int_0^{\infty}Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_8(a_8) = (\rho_0)_8(0)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - (f_{80})_8 a_8}$$

$$I_8(\gamma_8, a_8) = \int_0^a (\rho_0)_8(\tau_8)e^{\gamma_8\tau_8 + (f_{80})_8\tau + \int_0^\tau \mu_8(s_8)ds_8}d\tau$$

$$= \frac{(\rho_0)_8(0)}{\gamma_8}(e^{\gamma_8 a_8} - 1)$$

Let

$$E_8(a_8) = e^{-(f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau_8}$$

Then $[Y_8]_8[\gamma_8]_8, a_8) = E_8(a_8)e^{-\gamma_8 a_8}$. by a simple calculation,

$$\int_0^\infty \beta_8(a_8)Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma}\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}\ da\right);$$

$$\int_0^\infty Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma_8}\left(\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}\ da_8\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_8(0)}{\gamma_8}[f_{80}^8]\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}da_8\ (\int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8)$$

$$+\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right).$$

$$\left(1+\frac{(\rho_0)_8(0)}{\gamma_8}(f_{80}^{'})_8\left(\int_0^{\infty}E_8(a_8)da_8-\int_0^{\infty}E_8(a_8)e^{-\gamma_8 a_8}da_8\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right)(1+\frac{(f_{80}^{'})_8}{\gamma_8}(\rho_0)_8(0)\int_0^{\infty}E_8(a_8)da_8)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)da_8=1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\neq 0,$$

Since $\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)\int_0^{a}\alpha_8(\tau_8)Y^{-8}(\gamma_8,\tau_8)d\tau_8 da_8>0$. Therefore, the characteristic solution is

$$\gamma_8=-(f_{80}^{'})_8(\rho_0)_8\int_0^{\infty}E_8(a_8)da_8.$$

Since $f_8^{'}(P_0)_8>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_8),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_8),0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_8 + (\mu_8(a_8) + [f_8]_8(P_8^*) + [f_8]_8(Q_8^*))P_8^* = 0, \quad (26)$$

$$p_8^*(0) = h_8(Q_8^*)\int_0^\infty \beta_8(a_8)p_8^*(a_J)p_8^*(a_8)da_8 \quad (27)$$

$$g_8(P_8^*) - b_8 = 0, \quad (28)$$

Where $P_8^* = \int_0^\infty p_8^*(a_8)\, da_8$

Equation (28) yields $P^* = g^{-8}(b)$,

From (26),

$$p_8^*(a) = p_8^*(0)e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}$$

Substituting into (27), we need to solve

$$\left(1 - h_8(Q_8^*)\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}da_8\right)p_8^*(0) = 0$$

For $Q_8^*$

Fix $p_8^*(0)$ and $P_8^*$ Let

$$H_8(z_8) = 1 - h_8(z_8)\int_0^\infty \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}da_{,8}$$

And assume that

$$[k_3]_8 \int_0^{\infty} \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(g_8^{-8}(b_8))a_8 - [k_8]_8 a_8}\, da_8 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty} [H_8(z_8)] > 0.$$

Clearly,

$$\frac{dH_8}{dz_8} = h_8(z)(\,[f_8^{'}]_8(z_8)\,) \int_0^{\infty} a_8\beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - f_8(P_8^*)a_8 - f_8(z_8)a_8} da_8$$

$-h^{'}(z_8)\int_0^{\infty} \beta_8(a_8)\, e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}\ da_8$

By the assumptions on $\left[f_8^8\right]_8 and\ [h^8]^8 dH_8/dz > 0$. Hence, if

$$H_8(0) = 1 - \int_0^{\infty} \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8} da_8 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_8(0) < 0$ follows if $g_8^-(b_8) < [f_8^{-8}]_8(r_8)$, where $r_8$ is the solution of (12).

Indeed, since

$$R_8(s_8) = \int_0^{\infty} \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 e^{-y_8^a}} da_8$$

Is a decreasing function of $s_8$ when $P_8^* = [g^{-8}]_8(b_8) < [f_8^{-8}]_8(r_8)$, that is $[f_8]_8(P_8^*) < r$, $R\big(f_8(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_8(g^{-8}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8\left(g_8^{-8}(b)\right)a_8} > 1. \qquad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_8^*(a_8), Q_8^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_8^*(a_8), Q_8^*)$ is more complicated. We again consider perturbations $m_8(t_8 a_8) = p_8(t_8, a_8) - p_8^*(a_8)$ and $W_8(t) = Q_8(t) - Q^*$.Then

$$M_8(t) \equiv \int_0^\infty m_8(t, a_8)da_8 = P_8(t) - P_8^*,$$

$$[m_8]_8 + [m_a]_8 = [p_8]_8 + [p_a]_8 - \left[p^*\right]_8(a_8) \approx -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8$$

$$-[f_{8*}^8]_8 p_8^*(a_8)M - [f_{8*}^8]_8 p_8^*(\mathrm{a})W_8,$$

$$m_8(t,0)p_8(t,0) - p_8^*(0) \approx h_* \int_0^\infty \beta_8(a_8)m_8(t,a_8)\,da + \frac{[h_*^8]_8}{[h_*]_8}\, p_8^*(0)W_8$$

$$\frac{dW_8}{dt} = \frac{dQ_8}{dt} \approx [g_*^8]_8 Q_8^* M_{,8}$$

Where $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $\left[f_{8*}^{'}\right]_8$ denoting $\left[f^{'}{}_8(P^*)\right]_8\left[f_{8*}^{'}\right]_8$ denoting $[f_8]_8$ $(Q_8^*,$ $h_{*}^{'}$ denoting $h_*^{'}$ denoting $h_8^{'}$ $(Q_8^*)$, and $\left[g_*^{'}\right]_8$

Denoting $\left[\dot{g}\right]_8(P_8^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_8]_8 + [m_a]_8 = -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8 - \left[\dot{f}_{8*}\right]_8[p_8^*](a_8)M_8 - \left[\dot{f}_{8*}\right]_8 p_8^*(a_8)W_8$$

$$m_8(t,0) = [h_*]_8 \int_0^\infty \beta_8(a_8)m_8(t,a_8)da_8 + \frac{[h_*^8]_8}{[h^*]_8}[p^*]_8(0)\, W_{,8} \qquad (30)$$

$$\frac{dW_8}{dt} = [g_*^8]_8[Q_8^*]M_8$$

Integrating $[m_8]_8(t,a_8)$ with respect to a from 0 to ∞, we have

$$\frac{dM_8(t)}{dt} = \int_0^\infty [mt]_8(t,a_8)da_8$$

$$= -\int_0^\infty [m_a]_8\,(t,a_8)da_8 - \int_0^\infty \mu(a_8)m(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + \left[\dot{f}_{8*}\right]_8 \int_0^\infty p_8^*(a_8)da_8\right)M_8 - [f_{8*}^8]_8 \int_0^\infty p_8^*(a_8)da_8 W_8$$

$$= m_8(t,0) - \int_0^\infty \mu_8(a_8)\, m_8(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + \left[\dot{f}_{8*}\right]_8 P_8^*\right)M_8 - \left[\dot{f}_{8*}\right]_8 P_8^* W_{,8}$$

From (30), we arrive at the equations

$$\frac{dW_8(t)}{dt} = \int_0^\infty ([h_*]_8\beta_8(a_8) - \mu_8(a_8))m_8(t,a_8)da_8$$

$$-\left([f_{8*}] + [f_{8*}]_8 + [f_{8*}']_8 P_8^*\right) M_8 + \left(\frac{[h_*']_8}{[h_*]^8}\, P_8^*(0) - f_{8*}' P_8^*\right) W_8,$$

(31)

$$\frac{dW_8(t)}{dt}\left[g'_*\right]_8 Q_8^*\, M_{.8}$$

By the mean value theorem,

$$\int_0^\infty ([h_{"*}]_8 \beta_8(a_8) - \mu_8(a_8)) m_8(t, a_8) da_8 = ([h_*]_8 \beta_8(\theta_8) - \mu_8(\theta_8)) M_{,8}$$

Where $0< \theta_8 = \theta_8(t) < \infty$.

And if a $^8_8$ [a $^8_8$]2<4 [a $_8$]2<4 a $_8$e,

Re $\lambda_8 [\lambda_8]_8 = \text{Re}\lambda_8 [\lambda_8]_8 = -\frac{8}{8} a_8 [\text{a2}\ ]_8 < 0$.

Let ⊿=max$_{8<u<\infty}$ |h $_*$ [h $_*$] $_8$β $_8$($a_8$) − $\mu_8(a_8)$|. Then|$n_8$(t)| ⊿∀t . Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium(ρ $^*$ (a), $\boldsymbol{Q}^*$) is asymptotically stable if either a $^8_8$ <4 a $_8$ e and [a $_8$] > 1⊿,or [a $^8_8$] [a $_8$]e and (a $_8$ − $\sqrt{(a^8_8)_8 - 4(a_8)_8 e}$ > 1⊿

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+8}\right)$ and β $_8$($a_8$) = ae-$^{ma}$ Then ⊿=[$\mu_0$] + 2 Simple calculation leads to

$$\int_0^{\infty} \beta_8(a_8)\mathrm{e}^{-\int_0^{\mathrm{a}} \mu_8(\tau_8) + d\tau_8 \mathrm{da}} = \frac{1}{(\mathrm{m}+\mu_0+1)^8} + \frac{1}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_8 + [\mu_0]$ 0.5. Assume$[f_8](x_8) = K_8 x_8$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_8(x_8) = [0x_8]/(2 + x_8)$; the functional responses $[[f_8](x) = [0]\, x_8/(1 + x_8)$and $h_8(x_8)$=$(1+\omega_8 x_8)/(1 + x_8)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_8 + [\mu_0] + [\mathrm{K}][P_0])}{[m_8]_8 + (\mu_0)_8 + [\mathrm{K}][P_0])_8^3\ [\mathrm{K}]_8 [P_0]_8^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_8$i e $[P_0] < b/(\delta_8 - b_8)$then the equilibrium$(P_0)$ $[P_0]_8(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_8$

$[f_{8*}] = K_8\ \mathrm{P}^*\ [f_{8*}] = \frac{\sigma \mathrm{Q}_8^*}{8+\mathrm{Q}_8^*} \quad [h_*] = \frac{8+\omega_8 \mathrm{Q}_8^*}{8+\mathrm{Q}_8^*}$

$[f_{8*}] = K_8\ [\mathrm{f}_{8*}^8] = \frac{\sigma_8}{(8+\mathrm{Q}_8^*)^8}[\ \mathrm{g}_*^8] = \frac{\sigma_8}{(8+\mathrm{P}_8^*)^8} \quad [\mathrm{h}_*^8] = \frac{\omega_8 - 8}{(8+\mathrm{Q}_8^*)^8}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta$, and $\omega$ need to satisfy

$$\frac{\omega_8}{(m + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 + b_8} + \sigma)^8} + \frac{1\omega_8}{([m_8]_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 + b_8} + \sigma_8)^3} < 1$$

And

$$\frac{1}{(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^8} + \frac{1\omega_8}{(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8} + (m_8 + [\mu_0]_8 + \frac{K_8 b_8}{\delta_8 - b_8} + \sigma_8)\, Q_8^*)^3$

$= m_8 + [\mu_0]_8 + 3 + \frac{K_8 b_8}{\delta_8 - b_8}(1 + Q^*)(1 + \omega_8 Q_8^*$

Now

$$\left[a_{11}\right] - 4\left[a_{88}\right] e = \left(\frac{1K_8 b_8}{\delta_8 - b_8}\right)^8 + \frac{4K_8 b_8 \sigma_8 Q_8^*}{(\delta_8 - b_8)(8 + Q_8^*)} + \frac{\sigma_8^8 Q_8^{*8}}{(1 + Q_8^*)^8}$$

$$\frac{4\, b_8 \sigma_8 (\delta_8 - b) Q^*}{\delta_8 (1 + Q_8^*)^8} - \frac{4\, (\delta_8 - b_8)^8 (1 - \omega_8) Q_8^*}{\delta_8 (1 + \omega_8 Q_8^*)(1 + Q_8^*)} P^*(0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_8\, ae_8^{mu}\ [\mu_0]_8 - \frac{a_8}{a_8 + 8}$$

Hence if $[a_8^8 > 4(a_8)_8 e$ and

$$[a_8] = \frac{1K_8 b_8}{\delta_8 - b_8} + \frac{\sigma_8 Q_8^*}{1 + Q_8^*} > 1\Delta$$

**The nontrivial equilibrium is stable.**

IF we define

$$[a_8]_8 = [f_{8*}]_8 + [f_{8*}]_8 + [f_{8*}] P_8^*$$

$$[a_8]_8 = \left[f_{8*}^{8}\right] P_8^* - \frac{h_*^8}{h_*} P^*(0)$$

$$e = [g_*^8]_8 Q_8^*$$

$$\eta_8(t) = \left[h_*\right]_8 \beta_8 \left(\theta_8(t)\right) - \mu_8(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\left[a_8\right] M_8 - [a_8]_8 W_8 - [n(t)] M_8$$

$$w_8(t) = eM_8$$

Write

$$A = \begin{pmatrix} -[a_8]_8 - [a_8]_8 \\ e \qquad 0 \end{pmatrix}$$

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| $B=(t)=\begin{pmatrix} [\eta(t)]_8 & 0 \\ 0 & 0 \end{pmatrix}$ |
| And |
| $X(t)=\begin{pmatrix} M_8 \\ W_8 \end{pmatrix}$ |
| The solution of (32)is |
| $X_8(t) = e^{A,t}[X]_8(0) + \int_0 e^{At}$ |
| This equation notwithstanding its importance in all the models stand out for its importance for the following models<br><br>to which it is applied: |
| **Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016** |

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| (429)These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field. |
| (430)The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets. |
| (431) Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state. |
| (432) Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands. |
| (433)The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI |
| (434) The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9]. |
| (435)This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive. |

(436)More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

(437)Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.

(438)Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.

(439) The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.

(440)This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.

(441)The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.

(442)This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.

(443) For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.

(444)The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(445)Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.

(446)We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.

(447)When this is true—as is often the case—the system serves as (=) a "bath" for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called **thermalization.**

(448)However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.

(449)Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1).

(450)Bringing them to thermal equilibrium requires (e) an external bath.

(451) But so far, there has been no rigorous proof of (e) the existence of **such many-body localization**. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption.

(452)The concept of **localization** was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations).

(453)The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century.

(454)But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information.

(455)Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment.

(456)For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood.

(457) It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7].

(458) Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature.

(459) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.

(460)However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].

(461)Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.

(462) Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.

(463)As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked.

(464)His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice.

(465) An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site.

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| (466)This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself). |
| (467)This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition. |
| (468)However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime: |
| (469) These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects. |
| (470)It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies). |
| (471) This is a realistic scenario for (e) all known systems. |
| (472)On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects. |
| (473)On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other. |
| (474)No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true. |
| (475) His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found. |
| (476)Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization? |
| (477)Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects? |
| (478)Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions. |
| (479)So in those systems, the quantum phase transition does indeed exist. |
| (480)But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension. |
| (481)Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters. |

# PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION ATTRITION MODELS

## MODULE NINE

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Model is applied to following formulations Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society**

**Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice ($\alpha$=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP This paper is divided into four parts. In part A, some general considerations about (e&eb) gravitational radiation are followed by (e) a treatment of the scalar wave equation in the manner later to be applied to (e&eb) Einstein's field equations. In part B, a co-ordinate system is specified which is suitable for (e) investigation of outgoing gravitational waves from (e) an isolated axi-symmetric reflexion-symmetric system. The metric is expanded in (eb) negative powers of (e) a suitably defined radial co-ordinate r, and the vacuum field equations are investigated in detail. It is shown that the flow of information to infinity is controlled by (e&eb) a single function of (e) two variables called the news function. Together with initial conditions specified on a light cone, this function fully defines (eb) the behaviour of the system. No constraints of (e) any kind are encountered. In part C, the transformations leaving (eb) the metric in the chosen form are determined. An investigation of the corresponding transformations in (eb) Minkowski space suggests (eb) that no generality is lost by (e) assuming that the transformations, like the metric, may be expanded in (eb) negative powers of r. Gravitational Waves in General Relativity VII Waves from Axi-Symmetric Isolated Systems Royal Society Proceedings H. Bondi, M. G. J. van der Burg, A. W. K. Metzner Published 21 August 1962.DOI: 10.1098/rspa.1962.0161 STOP There is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant $\Lambda$, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive $\Lambda$ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the $\Lambda>0$ generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible. Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of (e) GW signals from BBHs with component spins $\chi_{1,2}\in[-0.99,0.99]$ aligned (e&eb) with the orbital angular momentum, component masses $m_{1,2}\in[2,48]M_\odot$, and (e&eb) total mass $M_{total}\leq 50\ M_\odot$. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin**

**bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins >+0.6 and component masses ≳20 M⊙. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201**

## 63. MODEL FORMULATION

Let $p_9(t, a_9)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(g_9)Q_9$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_9(0) = 0, \quad g_9(P_9) > 0, \ P_9 \ 0, \qquad \lim_{P_9 \to \infty} g_9(P_9) = (K_9)_9 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_9 + (\rho_a)_9 + D_9(a_9, P_9, Q_9)\rho_9 + 0, \qquad 0 < a_9 \ \infty,$$

$$\rho_9(t,0) = \int_0^\infty n_9(a_9, P_9, Q_9)\rho_9(t_9, a_9)da_9$$

Where $D_9$ is the mortality function and $n_9$ is the fertility function.

Put $_9$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_9}{dt}\right]_9 = x_9\left[\frac{dQ_9}{dt}\right]_9 - x_9\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_9}{dt}\right]_9 = x_3\left[\frac{dQ_3}{dt}\right]_9 - x_4\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_9}{dt}\right]_9 - x_6\left[\frac{dQ_3}{dt}\right]_9$$

Where $x_9, x_9, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_9$ and $g_9(P_9)Q_9$

Now when you add the term $(b_9Q_9)$ with $g_9(P_9)(Q_9)$ we are infect see is that $(b_9Q_9)$ is credited incrementally part of $g_9(P_9)(Q_9)$ or instant nearly, when subtracted from $\frac{dQ_9}{dt}$ , it means that $g_9(P_9)(Q_9)$ is slowly eaten by $\frac{dQ_9}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_9 = \mu_9(a_9) + (f_9)_9(P_9)_9 + (f_9)_9(Q)_9 \quad (2)$$

$$n_9 = \beta_9(a_9)h_9(Q_9), \quad (3)$$

Where $\mu_9$the age- dependent debit rate whit is $\mu_9(a_9)$ 0 $\forall a > 0$; $(f_9)_9$ is the density- dependent death rate with $f_9(0) = 0$, $(f_9^9)_9(P_9)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_9)_9(P_9) = (k_9)_9$ ; $f_9$ is the functional response of creditors to prey with $[f_9]_9(0) = 0, (f_9^9)_9(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_9)_9(Q_9) = [k_9]$ ; $\beta$ is the age dependent birth

rate with $\beta(a) \geq 0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h'(Q) \leq 0, \forall Q \geq 0$, and $\lim_{Q\to\infty} h(Q) = k_3 \geq 0$.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_9 + (\rho_a)_9 = -\mu_9(a_9) + (f_9)_9 P_9) + \big(f_9)_9(Q_9)\big)(\rho_9)_9 \tag{4}$$

$$\rho_9(t_9, 0) = h_9(Q_9) \int_0^\infty \beta_9(a_9)\rho_9(t, a_9)da, \tag{5}$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(P_9)Q_9 \tag{6}$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_9$ and $Q_9$ of (4-6) satisfy

$$0 \leq \rho_9(t, a_9) \leq \rho_9^*(t, a_9),$$

$$0 \leq Q_9(t) \leq Q_9^*(t),$$

Where $\rho_9^*(t, a_9)$ and $Q_9^*(t)$, are solutions of

$$(\rho_t^*)_9 + (\rho_a^*)_9 = -(\mu_9(a_9) + (f_9)_9(P_9^*) + (f_9)_9 Q_9^*))\, \rho_9^*, \tag{7}$$

$$\rho_9^*(t, 0) = \int_0^\infty \beta_9(a_9)\, \rho_9^*(t, a_9)da_9 \tag{8}$$

$$\frac{dQ_9^*}{dt} = -b_9 Q_9^* + g_9(P_9^*)Q_9^*, \tag{9}$$

With $P_9^*(t) = \int_0^\infty \rho_9^*\ (t, a_9)da_9$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_9^*(t)}{dt} = P_9^*(t)\big(r_9 - (f_9)(P_9^*) - (f_9)_9(Q^*)\big), \quad (10)$$

$$\frac{dQ_9^*(t)}{dt} = Q_9^*(t)(-b_9 + g_9(P_9^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_9(a)e^{-\int_0^a \mu_9(\tau)dt}e^{-ra}\ da_9 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_9^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(t)dt}\ da_9 > 1, \quad (H1)$$

Which guarantees $r_9 > 0$.

*Remark.* If $\beta_9(a) = [\beta_0]_9 a_9 e^{-\alpha_9 a_9}, \mu_9(a_9) \equiv \mu_9$ then the solution of (12) is

$$r_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_9]_9(P_9) = 0, [f_9](Q_9) = \lambda_9 Q_9\ h_9(Q_9) = 1, \quad and\ \ g_9(P_9) = c_9 P_9$$

$$[\rho_9]_9 + [\rho_a]_9 = -(\mu_9 + \lambda_9 Q_9)\rho_9$$

$$\rho_9(t,0) = \int_0^{\infty} [\beta_0]_9 a_9 e^{-\alpha_9 a_9} \rho_9(t,a_9)\, da_9$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + c_9 P_9 Q_9$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_9}{dt} = \omega_9 P_9 - \lambda_9 P_9 Q_9$$

$$\frac{dQ_9}{dt} = b_9 Q_9 - c_9 P_9 Q_9$$

Where $\omega_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$

THE EQUILIBRIUM $(\rho_0)_9(a_9), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_9(a_9), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_9 + (\rho_a)_9 = -\big(\mu_9(a_9) + (f_9)_9(P_9)\big)\rho,$$

$$\rho_9(t,0) = \int_0^{\infty} \beta_9\,(a_9)\rho_9(t,a_9)da_9 \qquad (13)$$

The solution of (13) has the form of

$$\rho_9(t,a_9) = Ce_9^{r(t-a_9)} - \int_0^a (\tau)d\tau - \int_0^9 (f_9)_9\big(P_9(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_9(t) = C_9 e^{r_9 t - \int_0^t (f_9)_9(P_9(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_9^{-9}(r)$ if $k_9 > r$. This is equivalent to

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{(k_9)_9 a_9} da_9 < 1.$$

In fact, if we define

$$R_9(s_9) = \int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau)d\tau} e^{-s_9 a_9}\, da_9$$

$R(s)$ is a decreasing function. $R_9(r_9) = 1$ And $R_9(k_9)_9 < 1$ imply $r_9 < (k_9)_9$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_9(a_9), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{k_9 a_9} da_9 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_9 \equiv \int_0^\infty (\rho_0)_{99}(a_9)da_9$ *satisfies*

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{-(f_9)_9(P_0)a_9} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_9(a_9), 0)$

Let $u_9(t,a_9) = \rho(t,a_9) - \rho_0(a_9)$ Then

$$N_9(t) \equiv \int_0^\infty u_9\,(t, a_9)da_9 = \int_0^\infty \rho_9(t, a_9)da_9 - \int_0^\infty (\rho_0)_9(a_9)da = P_9(t) - P_0)_9$$

And

$$(u_9)_9 + (u_a)_9 = (\rho_t)_9 + (\rho_a)_9 - (\rho_0^9)_9(a_9)$$

$$= -\mu_9(a_9) + (f_9)_9(P_9) + (f_9)_9(Q_9))\,(u_9 + (\rho_0)_9) + \mu_9(a_9)(\rho_0)_9 + (f_9)_9(P_0)_9(\rho_0)_9$$

$$= -\mu_9(a_9)u_9 - (f_9)_9(P_9) + (f_9)_9(Q_9))\,u_9 - (\rho_0)_9\,(f_9)_9(P_9) - (f_9)_9(P_0)_9$$

$$+ (f_9)_9(Q_9)), \qquad (15)$$

Where $(\rho_0^9)_9(a_9)$ denotes $d(\rho_0)_9(a_9)/da_9$

Expansion $(f_9)_9(P_9)$ and $(f_9)_9(Q_9)$ at $(P_0)_9$ and 0 yield respectively

$$(f_9)_9(P_9) = (f_9)_9((P_0)_9 + N_9)$$

$$= (f_9)_9(P_0)_9 + (f_9^9)_9(P_0)_9N_9 + o_9(N_9), \qquad (16)$$

$$(f_9)_9(Q_9)$$

$$= (f_9^9)_9(0)Q + o_9(Q_9) \qquad (17)$$

Write $f_9(P_0) = f_{90}, f_9'(P_0) = f_{90}'$, and $f_9'(0) = f_{90}'$ and substitute (16) and (17) into (15) to obtain

$$(u_9)_9 + (u_a)_9 \approx -(\mu_9(a_9) + (f_{90})_9)u - (f_{90}^9)_9(\rho_0)_9(N_9) - (f_{90}^9)_9(\rho_0)_9(Q_9)$$

$$u_9(t, 0) = \rho_9(t, 0)(\rho_0)_9(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t, a) + (\rho_0)_9(a_9)da_9 - (\rho_0)_9\,(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t, a_9)da_9 + (h_9(Q_9) - 1)(\rho_0)_9\,(0) \qquad (18)$$

Expanding

$$h_9(Q_9) = 1 + (h^9)_9(0)Q_9 + 0(Q_9) \equiv 1 + h_0'Q_9 + 0Q_9$$

And substituting it into (18), we arrive at

$$u_9(t,0) \approx \int_0^\infty \beta_9(a_9)u_9(t,a_9)da_9 + (h_{a_9}')(\rho_0)_9Q_9$$

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PREDATOR-PREY POPULATION MODELS

The linearization is the following :

$$[u_9] + [u_9] = -(\mu_9(a_9) + [t_{90}]_9u_9 - [\alpha_9]_9(a_9)N_9 - [\alpha_9]_9(a_9)Q_9$$

$$u_9(t,0) = \int_0^\infty \beta_9(a_9)u_9(t,a_9)da_9 + [\xi_9]Q_9 \qquad (19)$$

$$\frac{dQ_9}{dt} = -b_9Q_9$$

where

$$[\alpha_9]_9(a_9) = [t_9^9a]_9[p,o]_9(a_9)d2(a_9) = [f_{90}^9]_9(\rho\, o\,)_9(a_9)\xi = [\mathrm{h\ o}^9]_9[\rho\mathrm{o}]_9(\mathrm{o}) \text{ and } \mathrm{b}_9 = \mathrm{b}_9 - g_9(P_9)$$

If $g_9(P_0) > b$, clearly, $(\rho_0)_9(a_9), 0)$ is unstable and the predator population is persistent. Assume $g_9(P_0)_9 < b$.

Then, $Q_9(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_9 + u_a = -(\mu_9(a_9) + (f_{90})_9)u_9 - (\alpha_9)_9(a_9)N_9$$

$$= -(\mu_9(a_9) + f_{90})u_9 - \alpha_9(a_9)\int_0^\infty u_9(t,a_9)da_9$$

$$u_9(t,0) = \int_0^\infty \beta_9\,(a_9)u_9(t,a_9)da_9$$

Using the method developed by Cushing [2, 3], we suppose

$$u_9(t,a_9) = \tilde{u}_9(a_9)^{e^{\gamma(t-a_9)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_9}\,da_9$. Then,

A solution of the first differential equation for $\tilde{u}(a_9)$ gives,

$$\tilde{u}_9(a_9) = e^{-\int_0^a \mu_9(\tau_9)d\tau-(f_{90})_9 a}\left(\tilde{u}_9(0) - w\int_0^a e^{\int_0^\tau \mu_9(s_9)ds+(f_{90})}\tau_9\,(\alpha_9)_9(\tau_9)e^{\gamma_9\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\right)\tilde{u}_9(0)$$

$$+w\int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\int_0^a (\alpha_9)_9(\tau_9)e^{-\gamma_9\tau_9}e^{(f_{90})_9\tau_9+\int_0^\tau \mu_9(s_9)ds}d\tau_9\,da_9$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}da\right)(\tilde{u}_9(0))_9$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\int_0^a e^{(f_{90})_9\tau_9+\int_0^\tau \mu_9(s_9)ds+y\tau}\alpha_9(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(t_9)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_9\,(a_9)Y_9(\gamma_9,a_9)da_9\right)u_9(0)$$

$$+w\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)\int_0^a(\alpha_9)_9(\tau)_9(Y^{-9})_9(\gamma_9,\tau_9)d\tau_9\,da_9=0,$$

(24)

$$\left(\int_0^{\infty}Y_9(\gamma_9,a_9)da_9\right)\tilde{u}_9(0)$$

$$+\left(-1-\int_0^{\infty}Y_9(\gamma_9,a_9)\int_0^a(\alpha_9)_9(\tau)_9(Y^{-9})_9(\gamma_9,\tau_9)d\tau_9\,da_9\right)w=0.$$

Equation (20) has a solution of the form of $u_9(t,a_9)=u_9(a_9)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_9(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_9(\gamma_9,a_9)=\int_0^a(\rho_0)_9(\tau_9)Y^{-9}\,(\gamma_9,a_9)d\tau_9$. We obtain the characteristic equation

$$f_{90_9}^{l}\left(\int_0^{\infty}Y_9(\gamma_9,a_9)da_9\right)\left(\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)I_9(\gamma_9,a_9)da_9\right)$$

$$+\left(1-\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)da\right)\left(1+(f_{90}^{l})_9\int_0^{\infty}Y_9(\gamma_9,a_9)I_9(\gamma_9,a_9)da_9\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_9(a_9) = (\rho_0)_9(0)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - (f_{90})_9 a_9}$$

$$I_9(\gamma_9, a_9) = \int_0^a (\rho_0)_9(\tau_9)e^{\gamma_9\tau_9 + (f_{90})_9\tau + \int_0^\tau \mu_9(s_9)ds_9}d\tau$$

$$= \frac{(\rho_0)_9(0)}{\gamma_9}(e^{\gamma_9 a_9} - 1)$$

Let

$$E_9(a_9) = e^{-(f_{90})_9 a_9 - \int_0^a \mu_9(\tau_9)d\tau_9}$$

Then $[Y_9]_9[\gamma_9]_9, a_9) = E_9(a_9)e^{-\gamma_9 a_9}$. by a simple calculation,

$$\int_0^\infty \beta_9(a_9)Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma}\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}\ da\right);$$

$$\int_0^\infty Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma_9}\left(\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}\ da_9\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_9(0)}{\gamma_9}[f_{90}^9]\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}da_9\ (\int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9)$$

$$+\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right).$$

$$\left(1+\frac{(\rho_0)_9(0)}{\gamma_9}(f_{90}^{'})_9\left(\int_0^{\infty}E_9(a_9)da_9-\int_0^{\infty}E_9(a_9)e^{-\gamma_9 a_9}da_9\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right)(1+\frac{(f_{90}^{'})_9}{\gamma_9}(\rho_0)_9(0)\int_0^{\infty}E_9(a_9)da_9)=0.$$

We need ũ(a) $\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (94), to

$$1-\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)da_9=1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\neq 0,$$

Since $\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)\int_0^{a}\alpha_9(\tau_9)Y^{-9}(\gamma_9,\tau_9)d\tau_9 da_9>0$. Therefore, the characteristic solution is

$$\gamma_9=-(f_{90}^{'})_9(\rho_0)_9\int_0^{\infty}E_9(a_9)da_9.$$

Since $f_9'(P_0)_9>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_9),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_9),0)$ which is related to the extinction of predators (debit summation types: that which acts as a resource)and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $( p^*(a), Q^* )$. It must satisfy the equations:

$$[p_a^*]_9 + (\mu_9(a_9) + [f_9]_9( P_9^*) + [f_9]_9(Q_9^*) )P_9^* = 0, \quad (26)$$

$$p_9^*(0) = h_9(Q_9^*) \int_0^\infty \beta_9(a_9)p_9^*(a_J)p_9^*(a_9)da_9 \quad (27)$$

$$g_9(P_9^*) - b_9 = 0, \quad (28)$$

Where $P_9^*=\int_0^\infty p_9^*(a_9)\, da_9$

Equation (28) yields $P^* = g^{-9}(b)$,

From (26),

$$p_9^*(a) = p_9^*(0)e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}$$

Substituting into (27), we need to solve

$$\left(1 - h_9(Q_9^*)\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}da_9\right)p_9^*(0) = 0$$

For $Q_9^*$

Fix $p_9^*(0)$ and $P_9^*$ Let

$$H_9(z_9) = 1 - h_9(z_9)\int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9}\, da_{,9}$$

And assume that

$$[k_3]_9 \int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(g_9^{-9}(b_9))a_9 - [k_9]_9 a_9}\, da_9 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_9(z_9)] > 0.$$

Clearly,

$$\frac{dH_9}{dz_9} = h_9(z)\left([f_9']_9(z_9)\right)\int_0^\infty a_9\beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - f_9(P_9^*)a_9 - f_9(z_9)a_9} da_{\square}$$

$-h'(z_9)\int_0^\infty \beta_9(a_9)\, e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9}\ da_9$

y the assumptions on $\left[f_9^9\right]_9 and\ [h^9]^9 dH_9/dz > 0$. Hence, if

$$H_9(0) = 1 - \int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9} da_9 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_9(0) < 0$ follows if $g_9^-(b_9) < [f_9^{-9}]_9(r_9)$, where $r_9$ is the solution of (12).

Indeed, since

$$R_9(s_9) = \int_0^{\infty} \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 e^{-y\frac{a}{9}} da_9}$$

Is a decreasing function of $s_9$ when $P_9^* = [g^{-9}]_9(b_9) < [f_9^{-9}]_9(r_9)$, that is $[f_9]_9(P_9^*) < r$, $R\big(f_9(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_9(g^{-9}(b))) > 1, i.e.,$

$\int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9\left(g_9^{-9}(b)\right)a_9} > 1.$ (H4)

Then the system (4)-(6) has a unique positive equilibrium $(P_9^*(a_9), Q_9^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_9^*(a_9), Q_9^*)$ is more complicated. We again consider perturbations $m_9(t_9 a_9) = p_9(t_9, a_9) - p_9^*(a_9)$ and $W_9(t) = Q_9(t) - Q^*$.Then

$M_9(t) \equiv \int_0^{\infty} m_9(t, a_9) da_9 = P_9(t) - P_9^*,$

$$[m_9]_9 + [m_a]_9 = [p_9]_9 + [p_a]_9 - \left[p^*\right]_9(a_9) \approx -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9$$

$$-[f_{9*}^9]_9 p_9^*(a_9)M - [f_{9*}^9]_9 p_9^*(\mathrm{a})W_9,$$

$$m_9(t,0)p_9(t,0) - p_9^*(0) \approx h_* \int_0^{\infty} \beta_9(a_9) m_9(t, a_9)\, da + \frac{[h_*^9]_9}{[h_*]_9}\, p_9^*(0)W_9$$

$$\frac{dW_9}{dt} = \frac{dQ_9}{dt} \approx [g_*^9]_9 Q_9^* M_{,9}$$

Where $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $\left[f_{9*}^{'}\right]_9$ denoting $\left[f^{'}{}_9(P^*)\right]_9 \left[f_{9*}^{'}\right]_9$ denoting

$[f_9]_9$ $(Q_9^*,$ $h_9^*$ denoting $h_*^{'}$ denoting $h_9^{'}$ $(Q_9^*)$, and $\left[g_*^{'}\right]_9$

Denoting $\left[g^{'}\right]_9(P_9^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_9]_9 + [m_a]_9 = -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9 - \left[f^{'}{}_{9*}\right]_9 [p_9^*](a_9)M_9 - \left[f_{9*}^{'}\right]_9 p_9^*(a_9)W_9$$

$$m_9(t,0) = [h_*]_9 \int_0^\infty \beta_9(a_9)m_9(t,a_9)da_9 + \frac{[h_*^9]_9}{[h^*]_9}[p^*]_9(0)\, W_{,9} \qquad (30)$$

$$\frac{dW_9}{dt} = [g_*^9]_9[Q_9^*]M_9$$

Integrating $[m_9]_9(t,a_9)$ with respect to a from 0 to $\infty$, we have

$$\frac{dM_9(t)}{dt} = \int_0^\infty [mt]_9(t,a_9)da_9$$

$$= -\int_0^\infty [m_a]_9\,(t,a_9)da_9 - \int_0^\infty \mu(a_9)m(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + \left[f^{'}{}_{9*}\right]_9 \int_0^\infty p_9^*(a_9)da_9\right)M_9 - [f_{9*}^9]_9 \int_0^\infty p_9^*(a_9)da_9 W_9$$

$$= m_9(t,0) - \int_0^\infty \mu_9(a_9)\, m_9(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + \left[f_{9*}^{'}\right]_9 P_9^*\right)M_9 - \left[f_{9*}^{'}\right]_9 P_9^* W_{.9}$$

From (30), we arrive at the equations

$$\frac{dW_9(t)}{dt} = \int_0^{\infty} ([h_*]_9\beta_9(a_9) - \mu_9(a_9))m_9(t,a_9)da_9$$

$$-\left([f_{9*}] + [f_{9*}]_9 + \left[f_{9*}^{'}\right]_9 P_9^*\right) M_9 + \left(\frac{\left[h_*^{'}\right]_9}{[h_*]^9}\ P_9^*(0) - f_{9*}^{'} P_9^*\right) W_9,$$

(31)

$$\frac{dW_9(t)}{dt}\left[g_{*}^{'}\right]_9 Q_9^*\, M_{.9}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{"*}]_9\beta_9(a_9) - \mu_9(a_9))m_9(t,a_9)da_9 = ([h_*]_9\beta_9(\theta_9) - \mu_9(\theta_9))M_{,9}$$

Where $0< \ \theta_9 = \theta_9(t) < \infty$.

And if $a_9^9[a_9^9]2<4\ [a_9]2<4\ a_9e$,

$\text{Re}\,\lambda_9[\lambda_9]_9 = \text{Re}\lambda_9[\lambda_9]_9 = -\frac{9}{9}a_9[a2\ ]_9 < 0.$

Let $\Delta=\max_{9<u<\infty}\left|h_*[h_*]_9\beta_9(a_9) - \mu_9(a_9)\right|$. Then$|n_9(t)|$ ⍰⊿∀t ⍰0. Hence we obtain the following result.

**THEOREM 4**

Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_9^9 <4\ a_9e$ and $[a_9] > 1\Delta$,or $[a_9^9]$ ⍰4$[a_9]e$ and $(a_9 - \sqrt{(a_9^9)_9 - 4(a_9)_9e) > 1\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+9}\right)$ and $\beta_9(a_9) = ae^{-ma}$ Then $\Delta=[\mu_0] + 2$ Simple calculation leads to

$$\int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9) + d\tau_9 da} = \frac{1}{(m + \mu_0 + 1)^9} + \frac{1}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_9 + [\mu_0]$ 0.5. Assume$[f_9](x_9)= K_9 x_9$ Then the hypothesis (H9) is satisfied We take the conversion rate function $g_9(x_9)= [0x_9]/(2 + x_9)$; the functional responses $[[f_9](x) = [0]\, x_9/(1 + x_9)$and $h_9(x_9)=(1+\omega_9 x_9)/(1 + x_9)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_9 + [\mu_0] + [K][P_0])}{[m_9]_9 + (\mu_0)_9 + [K][P_0])_9^3\ [K]_9[P_0]_9^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_9$i e $[P_0] < b/(\delta_9 - b_9)$then the equilibrium$(P_0)$ $[P_0]_9(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_9$

$[f_{9*}] = K_9\ P^*\ [f_{9*}] = \frac{\sigma Q_9^*}{9+Q_9^*}$ $[h_*] = \frac{9+\omega_9 Q_9^*}{9+Q_9^*}$

$[f_{9*}] = K_9\ [f_{9*}^9] = \frac{\sigma_9}{(9+Q_9^*)^9}\ [g_*^9] = \frac{\sigma_9}{(9+P_9^*)^9}$ $[h_*^9] = \frac{\omega_9 - 9}{(9+Q_9^*)^9}$

In order to have (H3) and (H4) hold,$m, \mu_0, K, b, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_9}{(m+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9+b_9}+\sigma)^9}+\frac{1\omega_9}{([m_9]_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9+b_9}+\sigma_9)^3}<1$$

And

$$\frac{1}{(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9})^9}+\frac{1\omega_9}{(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9}+(m_9+[\mu_0]_9+\frac{K_9 b_9}{\delta_9-b_9}+\sigma_9)\,Q_9^*)^3$$

$$=m_9+[\mu_0]_9+3+\frac{K_9 b_9}{\delta_9-b_9}(1+Q^*)(1+\omega_9 Q_9^*$$

Now

$$\begin{bmatrix}1\\1\end{bmatrix}-4\begin{bmatrix}a_9\end{bmatrix}e=\left(\frac{1K_9 b_9}{\delta_9-b_9}\right)^9+\frac{4K_9 b_9\sigma_9 Q_9^*}{(\delta_9-b_9)(9+Q_9^*)}+\frac{\sigma_9^9 Q_9^{*9}}{(1+Q_9^*)^9}$$

$$\frac{4\,b_9\sigma_9(\delta_9-b)Q^*}{\delta_9(1+Q_9^*)^9}-\frac{4\,(\delta_9-b_9)^9(1-\omega_9)Q_9^*}{\delta_9(1+\omega_9 Q_9^*)(1+Q_9^*)}P^*(0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_9 \, ae^{mu}_9 \quad [\mu_0]_9 - \frac{a_9}{a_9+9}$$

Hence if $[a^9_9 > 4(a_9)_9 e$ and

$$[a_9] = \frac{1K_9 b_9}{\delta_9 - b_9} + \frac{\sigma_9 Q^*_9}{1 + Q^*_9} > 1\Delta$$

Some of the examples to which the model has been applied are:

**New look at the galaxy-shaping power of black holes**

(715) "It creates (eb) these bubbles, and the bubbles keep the plasma hot.

(716) That's what prevents galaxies from (e) becoming even bigger than they are now."

(717) Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.

(718) In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.

(719) Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.

(720) "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.

(721) "It's like nature's thermostat, which keeps these galaxies from (e) growing.

(722) If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."

**A new look at the galaxy-shaping power of black holes**

(723) An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the

(724) McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.

(725) Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.

**Doomed satellite missions**

(726) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.

(727) At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.

(728) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.

(729) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.

(730) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.

(731) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people

(732) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."

(733) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.

(734) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(735) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.

(736) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.

(737) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.

(738) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.

| |
|---|
| (739) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well. |
| (740) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma. |
| (741) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6. |
| (742) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus. |
| (743) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (744) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

| |
|---|
| (745) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (746) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (747) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (748) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (749) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (750) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (751) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (752) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (753) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (754) We know dark matter exists only from (e) its gravitational effects. |
| (755) It holds together galaxies and (e&eb) galaxy clusters. |
| (756) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (757) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |

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| (758) | They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (759) | X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (760) | They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (761) | This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (762) | An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |
| (763) | Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (764) | "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (765) | The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (766) | Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (767) | Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (768) | "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (769) | Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (770) | "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (771) | Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) |

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| consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories. |
| (772) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen. |
| (773) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer. |
| (774) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants. |
| (775) These were treated as (=) "aware" and "unaware" trials, respectively. |
| (776) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it. |
| (777) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another. |
| (778) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion. |
| (779) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness. |
| (780) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said. |
| (781) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness." |
| (782) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner. |
| (783) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain. |
| (784) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said. |
| (785) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

(786) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter.

(787) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined.

(788) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions.

(789) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.

(790) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.

(791) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.

(792) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.

(793) Noda and his colleagues instead studied the change in emissions over (e&eb) time.

(794) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.

(795) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).

(796) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.

(797) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.

(798) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson’s problem**

(799) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking “palm tree” organization.

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| (800) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society |
| (801) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible? |
| (802) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem. |
| (803) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other. |
| (804) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible. |
| (805) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges. |
| (806) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century. |
| (807) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching. |
| (808) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom. |
| (809) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth. |

**Spheres on trees**

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| (810) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration. |
| (811) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement. |
| (812) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum. |
| (813) By visualizing the problem in this way, the researchers noticed that these particular graphs don’t (e) have lots of branches extending from lots of (e) other branches. |
| (814) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure. |

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| (815) | The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network. |
| (816) | As a result, it provides (eb) an important clue for finding the global minimum. |
| (817) | It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum. |
| (818) | This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large. |

**Thompson's problem 2**

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| (819) | By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society |
| (820) | The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima. |
| **(821)** | The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node. |

**Implications for Thomson's and Smale's problems**

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| (822) | Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before. |
| (823) | "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around." |
| (824) | As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem. |
| (825) | "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view." |
| (826) | In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges. |
| (827) | From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear. |
| (828) | "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK. |
| (829) | "For larger numbers, there will likely be (=) multiple funnels. |
| (830) | Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum." |

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| (831) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem. |
| (832) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said. |
| (833) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org |

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| The nontrivial equilibrium is stable. |
| IF we define $[a_9]_9 = [f_{9*}]_9 + [f_{9*}]_9 + [f_{9*}]\,P^*_9$ |
| $$[a_9]_9 = \left[f^{9}_{9*}\right]P^*_9 - \frac{h^9_*}{h_*}P^*(0)$$ |
| $$e=[g^9_*]_9\,Q^*_9$$ |
| $$\eta_9(t) = \left[h_*\right]_9\beta_9\left(\theta_9(t)\right) - \mu_9\left(\theta(t)\right),$$ |
| The system(32)can be written as the following system: |
| $$[M(t)]\left[a_9\right]M_9 - [a_9]_9W_9 - [n(t)]M_9$$ |
| $$w_9(t) = eM_9$$ |
| Write |
| $$A=\begin{pmatrix} -[a_9]_9 - [a_9]_9 \\ e \qquad\quad 0 \end{pmatrix}$$ |

$$B=(t)=\begin{pmatrix} [\eta(t)]_9 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_9 \\ W_9 \end{pmatrix}$$

The solution of (32)is

$$X_9(t) = e^{A,t}[X]_9(0) + \int_0 e^{At}$$

**Note: There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and**

**sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by**

**term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation**

**alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation,**

**appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module**

**LEN NUNNEYS**

## Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

## Module numbered one

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems

and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1965); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be referred to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 6985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_6}{dt} = [N(t)]_6[(F_6)(R(t)]_6 - [M]_6[R(t)]_6$$

$$\frac{d[R(t)]_6}{dt} = [B]_6\,[(R)_6(t-T_6)] - [D]_6[(R(t)]_6 - [N(t)]_6\left[G\left(R(t)\right)_6\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_6}{dt} = a_6[N(t)]_6 - a_6[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_6 = a_3[N(t)]_3 - a_4[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_6 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

irst, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.6b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The

regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 6a is at the

**TIME DELAY**

FIG. 6. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.6$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**DELAY DUE TO PREDATOR MATURATION**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_6}{dt} = N_6(t - T_6)F_6(R_6(t - T_6) - N_6(t)M_6((R_6(t) \qquad (5)$$

$$\frac{dR_6(t)}{dt} = B_6\big(R_6(t) - D_6(R_6(t)\big) - N_6(t)G_6(R_6(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_6}{dR_6} + \frac{N_6 dG_6}{dR_6} - \frac{dB_6}{dR_6} > \frac{N_6 G_6}{2M_6}\left[\frac{dF_6}{dR_6} + \frac{dM_6}{dR_6}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_6/dR_6$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_6$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_6$) and the other derived from the non-delayed or "immediate" processes($I_6$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

ABSOLUTE STABILITY

$$L_6 = \frac{N_6 G_6}{2M_6}\frac{d(6nF_6)}{dR_6} (\geq 0) \qquad (6a)$$

$$I_6 = -\frac{N_6 G_6}{2M_6}\frac{d(6nM_6}{dR_6} (\geq 0) \qquad (6b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_6 > L_6 - I_6$, must be modified ahe $dM_6/dR_6$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_6 > L_6 - I_6, \quad \text{when} \quad P_6^2/4 > I \qquad (8a)$$

$$P_6 > L_6 - I_6 + \frac{(I_6 - P_6^2}{(L_6 - I)} \quad \text{when} \quad P_6^2/4 < I. \quad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

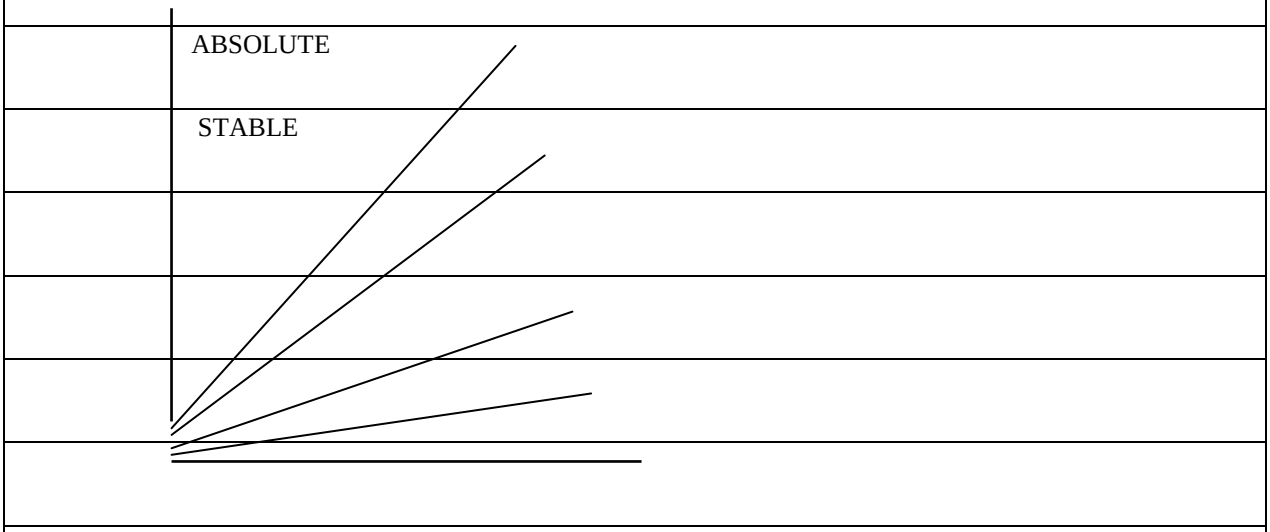



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

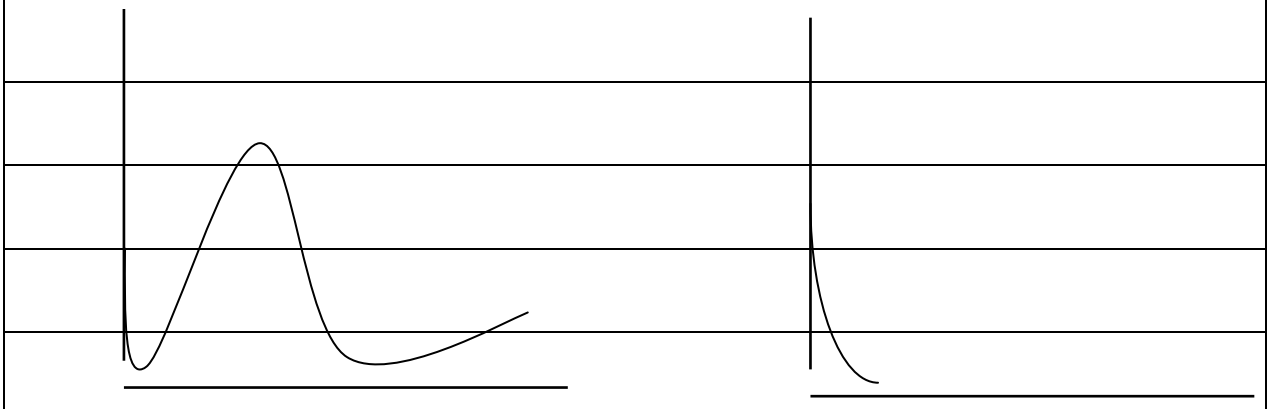
Can never completely offset the lagged effect,$L_6$; condition (8b) ensures that the total debit regulation, $P_6$, must always be positive even when $I_6$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_6$ and $L_6$. As $I_6$ increases relative to $L_6$ the zone of stability is enlarged; however,even when $I_6$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_6$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



**TIME DELAY**

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (6). Parameters in each case M=6; $dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6 = 0.05$ For (a) $NGdF/dR = 0.3, -N_6G_6dM_6/dR_6 = 0.6$; for (b)$NG_6dM_6/DR_6 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_6t}{dt} = N_6(t - T_6)F_6\big(R_6(t - T_6)\big) - N_6(t)M_6\big(R_6(t)\big)$$

(9)

$$\frac{dR_6t}{dt} = B_6\big(R_6(t - T_6)\big) - D_6\big(R_6(t)\big) - N_6(t)G_6\big(R_6(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) and the effective prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we

could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

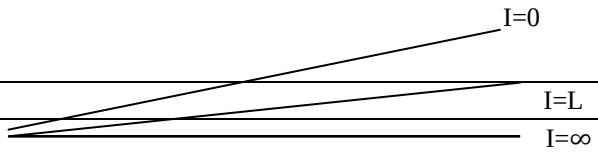
Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_6/dR_6$ is a measure of the control acting through the debit delay and setting $dB_6/dR_6 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_6/dR_6$ is a measure of the control acting through the credit delay although setting $dF_6/dR_6 = 0$ does not negate the effect of a predator delay; the value of $M_6$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_6/dR_6$and normalized by $M_6$ (see6a). The absolute stability of such a system is summarized in Fig. 6 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

TOTAL

ABSOLUTELY STABLE

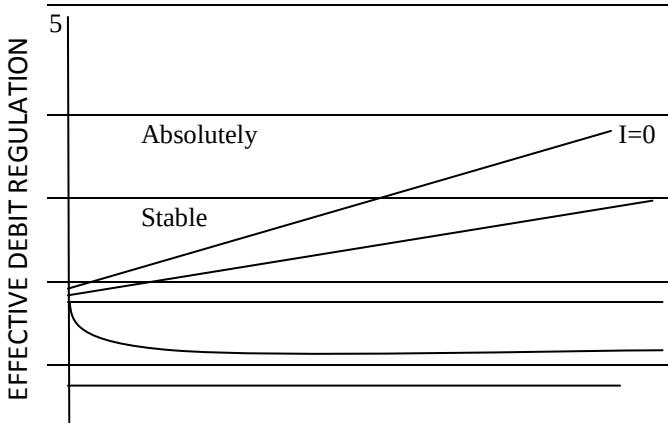



DELAYED DEBIT EFFECT

FIG6. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_6/dR_6 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.



-50 10

DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_6/dR_6 + N_6 dG_6/dR_6 + dB_6/dR_6)/M_6$). The immediate ($I_6$) and lagged ($L_6$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_6$(see 6b), is increased relative to the lagged part, $L_6$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_6 > I_6$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_6$is somewhat large than $I_6$ then the condition constraining prey regulation of recruitment ($dB_6/dR_6 < 0$) is

$$\frac{dD_6}{dR_6} + N_6\frac{dG_6}{dR_6} + \frac{dB_6}{dR_6} > \frac{(L_6 - I_6)M_6}{2} \qquad (10)$$

If $L_6 < I_6$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_6$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_6$

The - $D_6$ positions are defined by letting $\lambda_6 = iZ_6$ and separating (A1.2) into real and imaginary parts:

$$Z_6^2 + V_6 Z_6 \int_0^{\infty} \sin(Z_6 T_6) f(T_6) dT_6 - (X_6 + Y_6) = 0$$

$$Z_6 \left( U_6 - V_6 \int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6 \right) = 0$$

Noting that $X_6 + Y_6 = 0$ is we assume X+Y=0.It follows that $Z_6 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_6 = V_6 \int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6$$

If $U_6 > |V_6|$then given the constraints of $f(T_6)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A6.5)remains

greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

## ABSOLUTE STABILITY WITH A DELAY DUE TO CREDIT CONSOLIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_6^2 + (U_6 - V_6 + M_6)\lambda_6 - M_6\lambda_6 e^{-T_6\lambda_6} + [(U_6 - V_6)M_6 + Y_6] - [((U_6 - V_6 M_6 -]e^{-T_6\lambda_6} = 0$$

(For definitions see Appendix6) We follow the technique outlined in Appendix 6 and let $\lambda_6 = iZ_6$ and separate (A2.6) into real and imaginary parts

$$Z_6^2 + M_6 Z_6 \sin(Z_6 T_6) - (A_6 M_6 + Y_6) + (A_6 M_6 + X_6)\cos(Z_6 T_6) = 0$$

$$(A_6 M_6) Z_6 - M_6 Z_6 \cos(Z_6 T_6) + (A_6 M_6 - X_6)\sin(Z_6 T_6) = 0$$

Where $A_6 = U_6 - V_6$Using (A2.3), substitute for $Z_6^2$ and $Z_6$in (A2.2):

$$W_6^2(X_6^2 - 3M_6 A_6 X_6 - M_6^2(X_6 - Y_6) + 2M_6^3 A_6 + 2A_6^2 M_6^2]$$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_6 = 1 - \cos(Z_6 T_6)$.

In the absence of a delay the system is stable for $A_6 > 0,$ and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_6 > (X_6 - Y_6)/2\ M_6$ Thus a necessary condition for absolute stability is

$$A_6 > [\ X_6 - Y_6/2\ M_6$$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_6$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_6$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_6, A_6) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_6 \le 2$ *is* of interest First consider the extreme values

$$Q_6(0, A_6) = A_6^2\ (X_6 - Y_6)$$

$$Q_6(2, A_6) = (A_6 + 2\ M_6)^2\ (2A_6 M_6 - (X_6 - Y_6))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_6}{\partial A_6} = W_6^2\left(2M_6^3 + 4A_6 M_6^2 - 3M_6 X_6\right) + W_6\left(3M_6 A_6^2 - 2A_6 X_6 + 2M_6\left(2X_6 + Y_6\right)\right)$$

$$+2A_6\left(X_6+Y_6\right)$$

$=2M_6^2W_6^2(M_6+\mathrm{A}_6)+M_6A_6^2W_6+\left(2A_6M_6-X_6+Y_6\right)(A_6W_6+M_6W_6^2$

$+(2-W_6)((2M_6X_6+M_6Y_6)W_6+A_6(X_6+Y_6)),$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_6$quadratic in $\mathrm{W}_6$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_6$has no minimum in the range 0≤ $W_6$ ≤2.Defining

$Q_6\left(W_6,A_6\right)=a_6W_6^2+c_6 where\ a_6\ b_6\ and\ c_6$ Are functions of $A_6$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_6≥b_6≥0$$

And

$$b_6^2≤4a_6c_6$$

$$Case(i):Y_6=0$$

$Let\ Y_6 = 0$ And try the substitution A= $X_6 = /2M_6\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_6 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_6 X_6 /$ $2M_6$) is positive for *W<2. Since* $Q_6$ increases with $A_6$(see A2.9), increasing $A_6$will generate no further solutions of $Q_6$=0. Hence the condition A>X/2M defines absolute stability when $Y_6 = 0$.

$Case\ (ii)$: $Y_6 = y_6\ X_6$ ,

$Let\ Y_6 = y_6\ X_6$And try the substitution $A_6 = X_6/(2 + A_6\ y_6)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_6/(2 + n_6 y_6)$ we have

$$W_6^2\ y_6\left[X_6 n_6 \frac{(1 + n_6 y_6)}{(2 + n_6 y_6)} + M_6^2\left(2 - n_6 + n_6 y_6\right)\right]$$

$$-W_6\ X_6\left[\frac{X_6(1 + n_6 y_6)}{M_6^2(2 + n_6 y_6)^2} + 2\ y_6\left(n_6 - 1\right) + \frac{X_6^2(6 + y_6)}{M_6^2(2 + n_6 y_6)} = 0.\ (A2.12)\right.$$

Since $n_6$ and $A_6$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_6$

Defining $E_6 = (2 + n_6 y_6), the\ condition (A2.66?) can\ be\ written\ as$

$[\frac{X_6(E_6 - 6)}{E_6^2}]^2[\frac{X_6^2}{M_6^2} - 4\ y_6\ E_6^2\frac{X_6}{M_6^2} + 4y_6(n_6 - 2)E_6^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in ($X_6 / W_6^2$) can be expressed as

$2E_6[y_6E_6+(y_6E_6(6+y_6)(E_6-n_6))^{6/2}]>\frac{X_6}{M_6^2}>2E_6[y_6E_6-y_6E_6(6+y_6)(E_6-n_6))^{6/2}]$

**ABSOLUTE STABILITY**

Hence if $X_6/M_6^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0$W_6$2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_6-6)}{E_6^2}\left(\frac{X_6}{M_6^2}\right)^2-2y_6\frac{(E_6+n_6^2y_6)}{E_6}\left(\frac{X_6}{M_6^2}\right)-4y_6(E_6-n_6)0$$

And has the solution

$$\frac{X_6}{M_6^2}>\frac{E_6}{(E_6-6)}[y_6(E_6+n_6^2y_6)+[y_6^2(E_6+n_6^2y_6)^2+4y_6(E_6-6)(E_6-n_6)]^{6/2}]$$

Given $E_6>n_6$this lower limit to $X_6/M_6^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_6/M_6^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_6$in the appropriate range and this limit is further lowered by (A2.64)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_6<n_6(ie.,2A_6M_6^2<X_6-Y_6$)then (A2.6)and

(A2.8)show that $Q_6$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_6}{M_6^2} > 2E_6[y_6E_6 - (y_6E_6(6 + y_6)(E_6 - n_6)]^{6/2}]$$

This equations throws light on the individual general ledger and Nature's general ledger as we call it : see various papers on the issue. powers, which are the energies of Matrka we have already discussed, are not the only obstacles the yogi must overcome. He must, for example, also resist the temptation to rest content with the miraculous yogic powers (siddhi) he acquires in the course of his spiritual develop- ment. Again to do this he must practice Yoga. Similarly, in order to pervade the Fourth State gradually through the other states in the manner proper to practice at the Individual level, the method is the same. He must practise the higher yoga of the Tantras which, turning his mind inwards and freeing it from discursive representations, allows him to penetrate into the Supreme Principle. 312 Once the yogi has attained this contemplative state, his main problem is to make it permanent. In the introverted state the gross external movement of the breath is suspended and with it the activity of the intellect, mind, individualised consciousness, powers of the senses and the ego. 313 When the yogi rises out of this state, he is liable to fall again into the lower order of creation generated by Maya if he does not maintain his awareness of the higher reality he has experienced and allows his awakened, illumined insight to be obscured by the dream-like vision of thought-constructs. 314 Naturally, the yogi must rise out of the introverted condition of suspension. It is inherent in the very nature of reality that it should move out of itself. 315 Pure, universal consciousness initially transforms itself into the vital breath 316 charged with the impression iyasana) of the power of awareness attained through introversion. By attending to the pulse (spanda) of the breath as it moves out of the absolute, the

yogi can develop an intuitive sense of the inherent unity of all he will perceive in the mental and physical spheres created by the outpouring of consciousness. In this way he realises that his own nature is everywhere present in all he perceives and that all things thus reside within him. Blessed with this insight his consciousness remains free and unlimited even at the individual level where the breath, mind, senses and body are active. If the yogi fails to do this, he finds himself once again beset by the strictures of his embodied existence and must, as before, try to pervade all his other states of consciousness with the aesthetic delight (rasa) and wonder (camatkara) of the Fourth State he experienced in contemplation. Again this means that he must strengthen his pure, empowered awareness that his universal nature manifests as all things. 317 In this way he discovers Siva's presence in every sphere of individualised consciousness ranging from the breath to outer objectivity. The yogi's mind then becomes tranquil and undistracted because wherever it may wander, the yogi perceives only Siva, his authentic nature. 318 Consciousness is thus freed of all external referents and the yogi's subjectivity is purified of all identification with the body or anything else that belongs to the objective sphere. The yogi then becomes detached from the opposites of pleasure and pain and is transcendentally free (kevalin). 319 yogis are again, however, liable to fall if he allows himself to get entangled in the play of opposites. This fall is more serious than the others because, although he is caught by the confining restrictions of individual- ised consciousness as before, he is now also affected by karma. Fleeing from pain in the pursuit of pleasure he is bound to act (karma) to minimise one and maximise the other and so is thrown down to the lowest level of embodied subjectivity (sakala). In order to regain his lost state, he must ascend gradually, by Siva's grace, from one order of subjectivity to the next and so free himself progressively of the limitations of the lower levels to gain the greater freedom and expansion of the higher. As he progresses, the objective sphere also evolves from the grossest perceptions of physical objects outside the lowest order to subjectivity, through to the subtler inner, mental

perceptions to finally reach the order of subjectivity that contains objectivity within itself and is free to externalise it at will. 320 The degree to which this process develops depends, as before, on the yogi's awareness of the Fourth State. In consonance with the general principle that the remedy should suit the defect, the yogi is instructed to seek this higher state of consciousness in the wonder (camatkara) or delight (ananda) he feels in moments of intense physical pleasure. At first he experiences this subtle consciousness for an instant in the subjective sphere. If he manages to catch hold of it, it becomes more intense as the cognitive and objective spheres are also gradually pervaded and vitalised by it. Occasions for this practice are, for example, the sense of satisfaction one feels after a good meal or the aesthetic delight one experiences when listening to good music or the pleasure of sexual union with the Tantric consort or even solitary sexual excitation. In these moments of delight the yogi can penetrate momentarily into his own authentic Siva-nature (sambhavavesa) through the empowered contact (saktasparsa) m he makes with it in the freedom of the pure subjectivity of the Fourth State. 322 If the yogi develops his awareness of this higher level of consciousness and maintains it, he eventually experiences it constantly. 323 Clearly, what prevents the yogi from attending to his state of consciousness rather than the circumstances which induce it is the craving for pleasure (abhila§a) born of ignorance — the source of every impurity which clouds consciousness. Craving directs the yogi's attention towards outer, worldly things and so he is caught in the net of thought- constructs. 324 To free himself of his worldly desires and reverse this binding extroversion, the yogi must eradicate its cause. To be freed of all the ups and downs of the path and no longer be tormented by the possibility of a fall, the yogi must see reality perfectly and completely. This insight is itself liberation and the moment it dawns the yogi is instantly freed. This sudden realisation is the goal of Tantric Yoga. E Tantra declares: "He who perceives reality directly, even for the brief moment it takes to blink, is liberated that very instant and never reborn again." 325 Although the yogi's body and mind

continue to function as before, they are like mere outer coverings 326 which contain, but do not obscure, the mighty, universal consciousness which operates through them. The yogi's body is (=) the universe, the senses (=) the energies that (eb, eb+)vitalise it, his mind (=) Mantra, the rhythm of his breath (eb) the pulse of time and his inner nature pure, dynamic consciousness. Raised above all practice, and hence all possibility of falling to lower levels, the yogi realises that he has always been free 327 and that his journey through (e&eb) the **dark land of Maya was nothing but a dream, a construct of his own imagination** The channels of the senses by which the yogi enjoys the pleasure of the objects of sense are those through which he fills the three worlds with the consciousness emitted from his own Heart. 78 **The Doctrine of Vibration T**he senses are outer, subsidiary cycles of conscious energy (anucakra), revolving around the core or primary cycle (mukhyacakra) of pure consciousness. They scintillate with the diversity of sensations and resonate in the field of their awareness like the feathers of a peacock unfurled and folded back in the ecstasy of its dance. 79 Ultimately it is Siva Himself Who, for His own pleasure, 80 extends and withdraws them, 81 not the individual soul. The outward movement of the senses is the creative flow of consciousness from the subject who, like an ocean, is the source of the waves of their activity and the final resting place of the flow of the rivers of sensations and thought-forms that unfold through the arising of the power of each of the senses. 82 The Stanzas on Vibration declare: Indeed the individual soul does not activate the impulse of the will [which directs the body's activity] by himself alone but through his contact with [his] own [inner] strength {bald) made in such a way that he identifies with it [thus acquiring its power]. 83 The close contact between the senses and universal consciousness vitalises them and sets them in motion. They share in the properties of consciousness, just as a heated iron ball can burn, cook or give out light. 84 Spanda doctrines does not agree with the Sarnkhya view that the individual soul (purusa) alone stimulates activity in the objective sphere (prakrti) by the contact he makes with it. The soul can do

this only because universal consciousness animates him. Moreover, consciousness, which has the power to impel this activity, cannot itself be inactive as the Sarnkhya maintains. The influence it has on lifeless matter implies an active exertion. 85 This activity is the subtle, inner outpouring of the Self (svatmocchalatta) within the Self. Outside the confines of time and space, in the centre of the circle of the senses, this inner activity initiates the expansion (prasara) of the flow of conscious energy. As the pure exertion (udyoga, udyama) of Being (bhava), it is the source of the initial unfolding (unmesa) of Becoming. 86 Ksemaraja explains: [By] 'exertion' (udyama) is [meant] the emergence of the supreme [level of] intuition (pratibha) which is the sudden outpouring of the [immanent and transcendental] unfolding of the reflective awareness of pure consciousness. 87 The direct, intuition (pratibha) of reality, free of thought-constructs, is the supreme power (parasakti) of consciousness. The yogi can gain Divine **Body and Sacred Circle of the Senses 153** this pure intuitive awareness by grasping the initial expansion (unmesa) of consciousness which engenders the emergence of thought and the activity of the senses. It can be discovered in the subtle Centre, between one thought and the next. The Stanzas on Vibration teach: The expansion of consciousness that takes place when one is engaged in a single thought should be known to be the source from whence another arises. One should experience that for oneself! 88 This does not mean that the unfolding expansion of consciousness {unmesa) is an intermittent, transitory occurrence. 89 As the source of all thought, including the notions of past, present and future, 90 it stands outside the confines of time. Consciousness is in a state of perpetual expansion. It is an ever developing process, a tiny part of which we observe as the flow of events around us and the thoughts and feelings within us. The totality of this process can never be grasped objectively. Objectivity is a state of fragmentation (bheda) in which discrete elements are discernible and mutually distinguished as individual, specifically definable entities. It is the realm within which thought and language operate. To grasp reality in its completeness, we must go

beyond the partial representations of thought and speech. To experience the primordial source and basis of all things, we must pierce through the outer periphery of thought and plunge into the Centre, to discover the instant in which thought, and with it, the sense of diversity, initially emerges. Ksemaraja cites a practice from the Vijnanabhairava through which the yogi can achieve this state of awareness: Checking [the movement of] attention (cit) once it has quit its object, it must not [be allowed] to move to any other; contemplation (bhavana) then blossoms forth by [experiencing] the state in the centre between them (tanmadhyabhava). 91 Left unattended and undisciplined, awareness spontaneously shifts from one object of thought to (e&eb) the next. The transition from one to the other entails a movement through (e&eb) a state of pure indeterminate (nirvikalpa) awareness divested of all thought-forms and perceptions. The yogi who succeeds in checking the movement of his attention, experiences within himself the subsidence of thought into the emptiness of the Centre which unfolds full of the powerful pulsation of consciousness within which all objectivity merges and becomes one 92 in the trans- cendental outpouring (lokottarollasa) of its aesthetic rapture. 93 turning in on himself, the yogi must exert his power of awareness to (e) discern the vitality of (e) his own blissful Spanda nature. Just as a distant object becomes (=) more clearly visible by (e) exerting an effort to (e) perceive it, in the same way, through (e&eb) practice, the yogi's own Spanda nature becomes (=) progressively more evident to him, 94 The effort the yogi exerts to free himself of ignorance is essentially one with the force of awareness that impels the senses and mind. This inner exertion is the basis of every means to liberation. Ultimately bestowing the bliss of self- realisation, it carries the yogi along the path to salvation with all its twists and turns like wind which blows a cloth in various directions. 95 As the basis and ultimate goal of all yogic practice, this exertion is the pure wonder (camatkdra) inspired by the realisation of the essential identity of consciousness (cit) and being (sat). 96 As such it is the divine radiance (sphuratta) of the Heart of one's own consciousness 97 where the powers of

will, knowledge and action merge in the harmony (samarasya) of bliss. 98 Thus the highest form of worship (varivasya) is the contemplation of the inner strength (bald) or universal exertion of one's own true nature. 99 "One should examine [that reality]," says Ksemaraja, "with the exertion which is Bhairava's nature and is true service (sevana) to the integral inner nature [of all things] and the unfolding of the activity of one's own vitality (ojas) which withdraws all duality [back into consciousness]." 100 Mindful of the eternal joy (ahlada) of his inner strength, the yogi must give up all desire for the fleeting moments of petty pleasure (ksanikasukhalava) ]0] he may glean from outer objects. He must even give up the desire to achieve enlightenment and a reality (tattva) which can never be realised through a personal act of will. 102 In short, he must die to himself in an act of profound faith (sraddha) and adoration (adara) to find rest in the supreme reality free from the binding dichotomy of means and objectives (upeyopayabhava) — the reality of his perfected effort. The yogi, now full of devotion for this reality unfolding within himself and the universe (antarmukhatattva), merges into it and, by achieving the pure consciousness and bliss (cidananda) of the power (bala) of his own nature, achieves oneness. 103 The yogi is firmly established in his true nature by recognising the creative unfolding (unmesa) of consciousness through which he is able to know all things, whether in the present, past or future. 104 Master of the vitality which gives life to his mind and body, he can strengthenit directly without need of food or drink, 105 or quit it at will to vivify Divine Body and Sacred Circle of the Senses 155 that of another. 106 Experiencing the eternal satisfaction of his own uncreated nature, 107 the yogi is freed of disease and the infirmities of old age. 108 Inspired by this vitality, he has the inner strength to overcome the severest hardships and obstacles in the way of his ultimate goal. 109 The Stanzas on Vibration explain: Lassitude ravages the body and this arises from ignorance but if this is eliminated by an expansion of consciousness (unmesa), how can [ignorance], deprived of its cause, continue to exist? 110 Doubt (sanka) is the source of every spiritual ailment. 111 A

man in doubt about his true nature, and the path to follow by which he may come to realise it, is constantly overcome by difficulties. 112 "Due to ignorance" teaches the Sarvacaratantra, "the deluded man is in doubt and thus [suffers the cycle] of birth and death." 113 Caught between conflicting alternatives, the mind shifts from one to the other becoming more entangled in its own thoughts as it does so. Doubt thus contracts consciousness 114 sullying it with the turbulence of confused thinking. 115 At once both the root of "the ancient tree of transmigration" and the first sprout of its seed, 116 doubts deprives man of the innate bliss (sahajananda) of his own nature. 117 Like a thief it steals away the wealth of true knowledge and reduces man to an imagined state of poverty in which he feels hemmed in on all sides by constraints and limitations. 118 Ignorance and its sister, doubt, are the essence of all impurities (mala) which sully consciousness. To the degree in which the yogi is cured of this ailment, namely, the lassitude (glani) of doubt, his true nature becomes manifest, just as raw gold when heated is freed from dross. When the yogi recognises the all-powerful expansion (unmesa) of his consciousness, he is projected beyond the realm of relative distinctions between virtue and vice. Thus overcoming all doubt as to what he should or should not do, he penetrates into the pulsing Heart of Bhairava's consciousness. 119 As the Stanzas on Vibration say: An individual who [though] desirous of doing various things [but] is incapable of doing them due to his innate impurity, [experiences] the Supreme state when the disruption (ksobha) [of his ego] ceases. Then [the soul realises] that his [true] uncreated nature is [universal] agency and perceiving subjectivity and so he knows and does whatever [he] desires. 120 the individual soul is(=) a combination of (e) conscious (cit) and 156 The Doctrine of Vibration unconscious (acii) elements 121 brought together when he allows the power of his authentic nature to be obscured by (e) his own impurity. This impurity is (=) three-fold. The first is technically called (=) the 'Impurity of Individuality' (anavamald). Due to (e) this impurity, the individual soul fails to (e) recognise his all-embracing fullness (purnata) and is disturbed by (e) his craving

for (e) experience (bhoga) based on (e) the unreasoned assumption that he is (=) incomplete (apurnd) and in need of (e) something outside himself. 122 Thus it manifests as (=) desire or attachment (rdga). ]23 At the same time it is (=) primordial ignorance, 'the cause of the sprout of (e) transmigration'. 124 As such, it works in two ways. Firstly, it deprives (e) individual consciousness of freedom. The yogi affected by (e&eb) this impurity is incapable of (e) entering or emerging out of the states of contemplation as he wishes. Secondly, this impurity renders (eb) the individual unconscious of his freedom, which for(e) the yogi means that he is (=) liable to lose consciousness of (e) his true nature when he is (=) not in a state of (e) contemplative absorption. 125 The second impurity is (=) called the 'Impurity of Maya' (maylyamald). This comes into operation when consciousness has been contracted by (e) the Impurity of Individuality. It limits (e) the power of knowledge operating in (e&eb) the individual soul through (e&eb) the inner and outer organs of (e) perception. 126 Due to this impurity he perceives (e) diversity everywhere. 127 Finally, the third impurity, that of Karma, comes into play when (e) the individual, deprived of (e) the freedom and knowledge of (e) unconditioned consciousness, acts (e&eb) in his limited way prompted by (e) desires and fears for (e) his personal gain. Disturbed by (e) these three impurities the soul cannot find rest in (eb) himself.

And

$$E_6 > n_6$$

Replacing $E_6$ and $n_6$ in these conditions leads to the form

If $A_6^2 > 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6$.

If $A_6^2 < 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6 .+ (Y_6 - A_6^2/2)^2/2M_6(X_6 - Y_6)$.

Note that by normalizing $A_6$ by $M_6$ and $X_6$ and $Y_6$ by $W_6^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_6^2 + (M_6 + U_6)\lambda_6 - (M_6 + V_6)\lambda_6 e^{2T_6\lambda_6} + (M_6 U_6 + Y_6)$$

$$-[M_6(U_6 + V_6) - X_6]\, e^{2T_6\lambda_6} + M_6 V_6\, e^{2T_6\lambda_6} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_6 = \mathrm{i}\, Z_6$ *and* separate

(A3.1)into real and imaginary parts

$$Z_6^2 + (M_6 + V_6) Z_6 \sin(Z_6 T_6) - (U_6 M_6 + Y_6) + (M_6(U_6 + V_6) - X_6)$$

$$\mathrm{X}\cos(Z_6 T_6) - 2M_6 V_6 \cos^2(Z_6 T_6) + M_6 V_6 = 0$$

$$(M_6 + U_6) Z_6 - (M_6 + V_6) Z_6 \cos(Z_6 T_6) + (M_6(U_6 + V_6) - X)$$

$$\mathrm{X}\sin((Z_6 T_6) - M_6 V_6 \sin(Z_6 T_6) = 0$$

Using (A.3.3) to eliminate $Z_6$ and $Z_6^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_6 =$

$1 - \cos Z_6 T_6$ and $A_6 = U_6 - V_6$ we obtain

$$4M_6^2 V_6^2 W_6^4 + 2M_6 V_6 W_6^3[(M_6 - V_6)^2 + A_6(M_6 + V_6) + 2(A_6 M_6 - X_6)]$$

$$+ W_6^2 [M_6 - V_6)^2 (2 A_6 M_6 + Y_6 - X_6) + (2A_6 M_6 - X_6)^2 - A_6 (M_6 - V_6)(2A_6 M_6 - X_6)$$

$$+ M_6 V_6 (A_6^2 + 4(X_6 + Y_6))]$$

$$+ W_6 [A_6^2 + 2X_6)(2 A_6 M_6 - X_6) + 2A_6 Y_6 (M_6 + V_6) - M_6 A_6^3] + A_6^2 (X_6 M_6 - Y_6) = 0$$

Being a quartic in $W_6$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_6 > 0$$

$$and$$

$$C_6 > (X_6 - Y_6)/2 M_6$$

$recalling\ that\ A_6 = U_6$ and defining $C_6 = U_6 + V_6$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_6 (0, U_6 V_6 = A_6^2 (X_6 + Y_6)$$

$$Q_6 (2, U_6 V_6) = (C_6 + 2M_6^2)^2 (2C_6 M_6 + Y_6 - X_6)$$

So that given (A3.6) the extreme $W_6$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_6$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_6$ the existence of a D-position at this point is

determined by letting cos $(Z_6 T_6) = -1\ in\ (A3.3)$

ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 e^{-\tau_6\lambda_6} + X_6 + Y_6 = 0 \qquad (A1.1)$$

Where $V_6 = dB_6/dR_6$ $U_6 = dD_6/dR_6 + N_6 dG_6/dR_6$ $X_6 = N_6G_6dF_6/dR_6$ and $Y_6 = -N_6G_6dM_6/dR_6$ (evaluated at the equilibrium). Note that $U_6X_6$ and $Y_6$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_6$ as $(T_6)$ . (It follows that $f(T_6)$ 0 for all $T$, and $\int_0^\infty f(T_6)dt = 1.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 \int_0^\infty e^{-\tau_6\lambda_6} f(T_6) + X_6 + Y_6 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MODULE NUMBERED TWO**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed

the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).. Haimovici's and Nunny's models have been applied to some of the Deleuzean and Heideggerian problems too : Scoring on that count are :

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time

delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t - T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if

only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

TIME DELAY

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

## MATURATION AND CONSUMMATION OF CONFEDERATION OF DEBIT AND CREDIT TYPES:

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t-T_7)F_7(R_7(t-T_7) - N_7(t)M_7((R_7(t) \qquad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7} + \frac{N_7 dG_7}{dR_7} - \frac{dB_7}{dR_7} > \frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7} + \frac{dM_7}{dR_7}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes $(L_7)$ and the other derived from the non-delayed or "immediate" processes$(I_7)$:as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_7 = \frac{N_7 G_7}{2M_7}\frac{d(7nF_7)}{dR_7} (\geq 0) \qquad (7a)$$

$$I_7 = -\frac{N_7 G_7}{2M_7}\frac{d(7nM_7}{dR_7} \ (\geq 0) \qquad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7 > L_7 - I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_7 > L_7 - I_7, \quad \text{when} \quad P_7^2/4 > I \qquad (8a)$$

$$P_7 > L_7 - I_7 + \frac{(I_7 - P_7^2}{(L_7 - I)} \quad \text{when} \quad P_7^2/4 < I. \qquad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

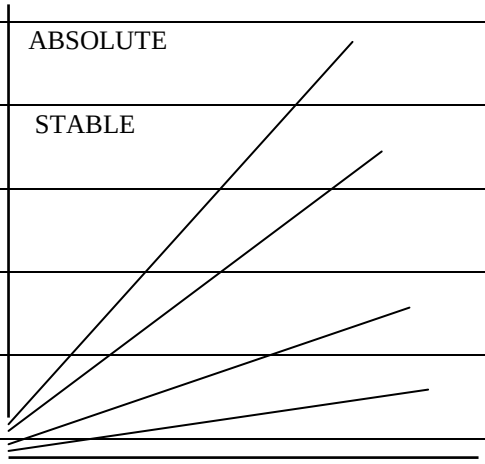



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit

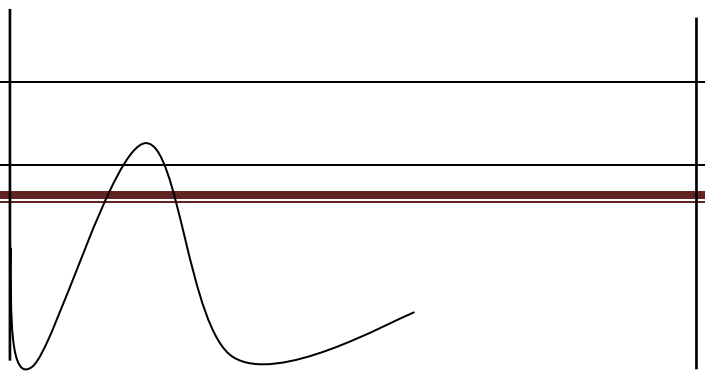
(like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

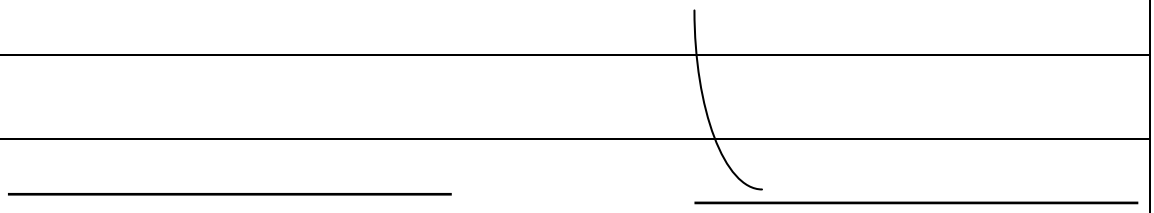
It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the





TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t-T_7)F_7\big(R_7(t-T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t-T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation $(dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7)$ and the effective prey regulation $(dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7)$ must be positive.

The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the

total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

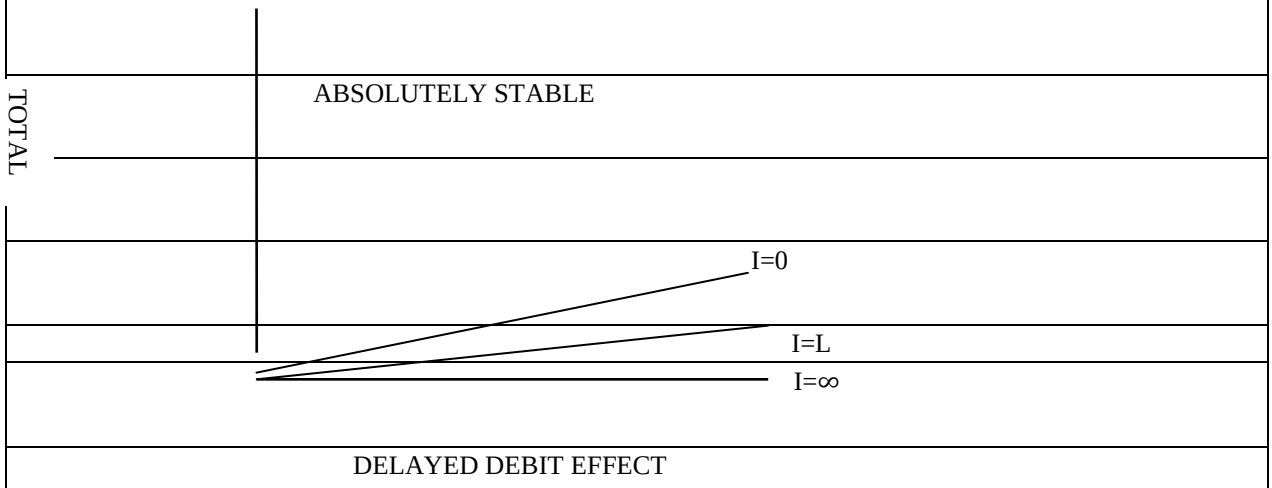


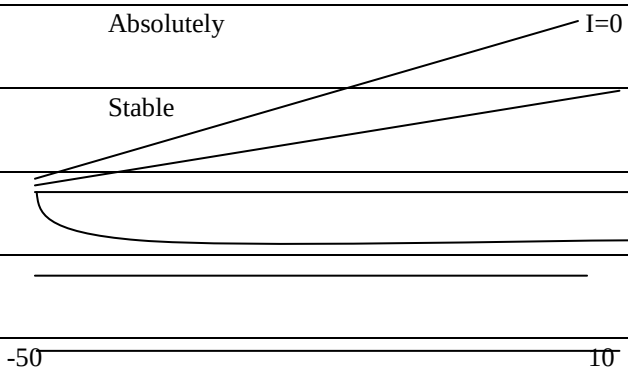

FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.






**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay.

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^\infty \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7\left(U_7 - V_7\int_0^\infty cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7\right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7\int_0^\infty cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$then given the constraints of $f(T_7)$, $(A1.5)$*can never be satisfied Co*nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7\lambda_7 e^{-T_7\lambda_7} + [(U_7 - V_7)M_7 + Y_7] - [((U_7 -$$

V 7M 7−]$e^{-T7\lambda 7}$=0

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$Z_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7)\cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7)\sin(Z_7 T_7) = 0$$

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $Z_7^2$ and $Z_7$in (A2.2):

$W_7^2(X_7^2 - 3M_7 A_7 X_7 - M_7^2(X_7 - Y_7) + 2M_7^3 A_7 + 2A_7^2 M_7^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_7$=1-cos $(Z_7 T_7)$.

In the absence of a delay the system is stable for $A_7 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7 > (X_7 - Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$$A_7 > [\ X_7 - Y_7/2\ M_7$$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_7, A_7) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_7 \le 2$ *is* of interest First consider the extreme values

$$Q_7(0, A_7) = A_7^2\ (X_7 - Y_7)$$

$$Q_7(2, A_7) = (A_7 + 2\,M_7)^2\,(2A_7 M_7 - (X_7 - Y_7))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7 M_7^2 - 3M_7 X_7\right) + W_7\left(3M_7 A_7^2 - 2A_7 X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7+Y_7\right)$$

$=2M_7^2W_7^2(M_7+\mathrm{A}_7)+M_7A_7^2W_7+\left(2A_7M_7-X_7+Y_7\right)(A_7W_7+M_7W_7^2$

$+(2-W_7)((2M_7X_7+M_7Y_7)W_7+A_7(X_7+Y_7)),$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $W_7$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$has no minimum in the range $0 \le W_7 \le 2$.Defining

$Q_7\left(W_7,A_7\right)=a_7W_7^2+c_7 where\ a_7\ b_7\ and\ c_7$ Are functions of $A_7$we can define the condition for a minimum in the range 0-2

condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7 \ge b_7 \ge 0$$

And

$$b_7^2 \le 4a_7c_7$$

$$Case(i): Y_7=0$$

$Let\ Y_7=0$ And try the substitution A= $X_7=/2M_7\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}\ (W_7 X_7 / 2M_7)$ is positive for *W<2. Since* $Q_7$ increases with $A_7$(see A2.9), increasing $A_7$will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

$Case\ (ii): Y_7 = y_7\ X_7$ ,

*Let* $Y_7 = y_7\ X_7$And try the substitution $A_7 = X_7/(2 + A_7\ y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + n_7 y_7)$ we have

$$W_7^2\ y_7\left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2\left(2 - n_7 + n_7 y_7\right)\right]$$

$$- W_7\ X_7\left[\frac{X_7(1 + n_7 y_7)}{M_7^2(2 + n_7 y_7)^2} + 2\ y_7\left(n_7 - 1\right) + \frac{X_7^2(7 + y_7)}{M_7^2(2 + n_7 y_7)}\right] = 0.\ (A2.12)$$

Since $n_7$ and $A_7$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7), the\ condition (A2.77?) can\ be\ written\ as$

$[\frac{X_7(E_7 - 7)}{E_7^2}]^2[\frac{X_7^2}{M_7^2} - 4\ y_7\ E_7^2\frac{X_7}{M_7^2} + 4y_7(n_7 - 2)E_7^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$2\ E_7[y_7 E_7 + (y_7\ E_7(7 + y_7)(E_7 - n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - y_7 E_7(7 + y_7)(E_7 - n_7))^{7/2}]$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 W 7 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7 - 7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2 - 2y_7\frac{(E_7 + n_7^2 y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right) - 4y_7(E_7 - n_7)$$ 0

And has the solution

$\frac{X_7}{M_7^2} > \frac{E_7}{(E_7-7)}[y_7(E_7 + n_7^2 y_7) + [y_7^2 (E_7 + n_7^2 y_7)^2 + 4 y_7 (E_7 - 7)(E_7 - n_7)]^{7/2}]$

Given $E_7 > n_7$ this lower limit to $X_7 / M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7 / M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$ in the appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7 < n_7$ $(ie., 2A_7 M_7^2 < X_7 - Y_7)$ then (A2.7)and (A2.8)show that $Q_7$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - (y_7 E_7 (7 + y_7)(E_7 - n_7)]^{7/2}]$$

And

$$E_7 > n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2 > 2Y_7$then $A_7 > (X_7 - Y_7)/2M_7$.

If $A_7^2 < 2Y_7$then $A_7 > (X_7 - Y_7)/2M_7$.+ $(Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$and $X_7$and $Y_7$by $W_7^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSILIDATAN**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7\lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]e^{2T_7\lambda_7} + M_7 V_7 e^{2T_7\lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i}\, Z_7$ *and* separate

(A3.1)into real and imaginary parts

$$Z_7^2 + (M_7 + V_7)Z_7 \sin(Z_7 T_7) - (U_7 M_7 + Y_7) + (M_7(U_7 + V_7) - X_7)$$

$$\mathrm{X}\cos(Z_7 T_7) - 2M_7 V_7 \cos^2(Z_7 T_7) + M_7 V_7 = 0$$

$$(M_7 + U_7)Z_7 - (M_7 + V_7)Z_7 \cos(Z_7 T_7) + (M_7(U_7 + V_7) - X)$$

$$X \sin\left(\left(Z_7 T_7\right) - M_7 V_7 \sin\left(Z_7 T_7\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_7$ and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - cos\ Z_7\ T_7$ and $A_7 = U_7 - V_7$ we obtain

$$4 M_7^2 V_7^2 W_7^4 + 2 M_7 V_7\ W_7^3 [(M_7 - V_7)^2 + A_7 (M_7 + V_7) + 2(A_7 M_7 - X_7)]$$

$$+ W_7^2 [M_7 - V_7)^2 (2 A_7\ M_7 + Y_7 - X_7) + (2A_7\ M_7 - X_7)^2 - A_7 (M_7 - V_7)(2A_7\ M_7 - X_7)$$

$$+ M_7 V_7 (A_7^2 + 4(X_7 + Y_7))]$$

$$+ W_7 [A_7^2 + 2X_7)(2 A_7\ M_7 - X_7) + 2A_7 Y_7 (M_7 + V_7) - M_7 A_7^3] + A_7^2 (X_7\ M_7 - Y_7) = 0$$

Being a quartic in $W_7$ (A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7 > 0$$

$$and$$

$$C_7 > (X_7 - Y_7)/2\ M_7$$

$recalling\ that\ A_7 = U_7$ and defining $C_7 = U_7 + V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0, U_7 V_7 = A_7^2 (X_7 + Y_7)$$

$$Q_7\left(2, U_7 V_7\right) = \left(C_7 + 2M_7^2\right)^2 (2C_7 M_7 + Y_7 - X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1$ *in* $(A3.3)$

## ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

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# Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price, University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss**

(2329) In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them.

(2330) Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors.

**Classical vs. relativistic black holes**

(2331) Something like a black hole exists within (eb) Newton's classical theory of gravity.

(2332) In that theory, an energy argument tells us that there is (=) an escape velocity vesc=2GM/R−−−−−−−√ from (e) the surface of any spherical object of (e) mass M and radius R

(2333) If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity.

(2334) Thus the condition for such an "unseeable" object is R<2GM/c2

(2335) In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape.

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| (2336) This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity. |
| (2337) To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it. |
| (2338) The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c . |
| (2339) This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed. |
| (2340) An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity. |
| (2341) Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable. |

**Stationary black holes**

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| (2342) In the Newtonian theory, gravity is described by (e) the potential $\Phi$. |
| (2343) Inside a spherical object the form of (e) $\Phi(r)$ depends on (=) the interior structure, but in the vacuum outside (e) matter the potential $-GM/r$ depends only on (=) the interior mass. |
| (2344) Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object. |
| (2345) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward. |
| (2346) This horizon radius $RH=2GM/c2$ is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory. |
| (2347) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution. |
| (2348) Radius RHhere actually means that the area of the event horizon is (=) $4\pi R2H$. |
| (2349) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution. |
| (2350) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity. |
| (2351) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own. |
| (2352) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole. |
| (2353) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density. |
| (2354) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed. |
| (2355) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole. |

(2356) In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter.

(2357) Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum.

(2358) In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior.

(2359) The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object.

(2360) Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum.

(2361) These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.]

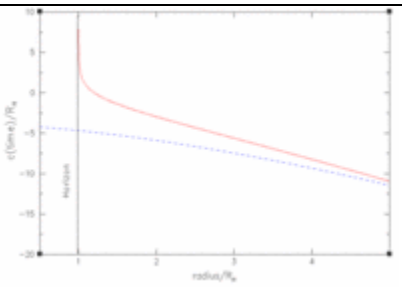
(2362) All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time.

(2363) But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?"

(2364) The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point.

(2365) One of these coordinates is (=) called the "coordinate time."

(2366) Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate.



(2367) Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole

(2368) At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole.

(2369) Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations.

(2370) For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree

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| (2371) An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events. |
| (2372) The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving. |
| (2373) Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle. |
| (2374) The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return. |
| (2375) But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon. |
| (2376) Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future. |
| (2377) The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?" |

**Black hole parameters**

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| (2378) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin). |
| (2379) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun. |
| (2380) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits. |
| (2381) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form. |
| (2382) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds. |
| (2383) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2. |
| (2384) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed. |
| (2385) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer. |
| (2386) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2) |
| (2387) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials. |
| (2388) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable |
| (2389) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite. |
| (2390) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole. |

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| (2391) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars. |
| (2392) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity. |
| (2393) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity. |
| (2394) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite. |
| (2395) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity. |
| (2396) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum. |
| (2397) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship. |
| (2398) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= (4πG2/c4)[(M+M2−a2−−−−−−−√ )2+a2]. |
| (2399) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored. |

**Dynamical black holes**

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| (2400) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape. |
| (2401) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, **a horizon can be born**. |
| (2402) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid. |
| (2403) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward. |
| (2404) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations. |
| (2405) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes. |
| (2406) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3. |
| (2407) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/ |
| (2408) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers. |
| (2409) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission. |

(2410) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.

(2411) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.

(2412) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.

(2413) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.

(2414) The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature $6\times10-8$ (Mo/M) K.

(2415) Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.

**Astrophysical black holes**

(2416) Black holes in our Universe can be grouped as: primordial black holes, stellar-mass black holes, and supermassive black holes. The first is highly speculative; the second and third are broadly accepted.

(2417) Primordial black holes of all masses are postulated to have formed from (e) quantum fluctuations in the early Universe, but those with (e&eb) mass less than around 1013kg would have already evaporated due to Hawking radiation.

(2418) No significant observational evidence has yet been found of (e) the existence of these objects [MacGibbon and Carr (1991)].

(2419) Stellar-mass black holes, ranging from a few to a few tens of solar masses, are (=) normal but rare endpoints in (eb) the evolution of massive stars.

(2420) When a star exhausts its nuclear fuel and cools, it must collapse unless (e) it is supported by nonthermal forces.

(2421) It is known that such nonthermal forces cannot (e) resist gravitational compression for (e) masses greater than around 1.5 to 3Mo. (The uncertainty is due to (e) uncertainty in our knowledge of nuclear physics at high densities.)

(2422) There are many stars far more massive than this, but in their death throes massive stars expel (eb) a great deal of their mass in supernova explosions

(2423) Even stars initially as massive as 20-30Mo may blow off enough mass to (e) leave a remnant neutron star smaller than 1.5Mo.

(2424) Stars more massive than 20-30Mo may form (eb) a neutron star core, only to have it collapse due to (e) fallback of material from (e) the stellar mantle. Still more massive stars (above ~40Mo) may form (eb) a black hole directly in collapse, with or without (e&eb) a supernova explosion.

(2425) These masses are (=) rather uncertain at present, as our understanding of supernova explosion mechanisms is (=) still evolving [Woosley and Janka (2005), Muno (2006)].

(2426) Collapsing gas clouds larger than ~100Mo are expected to dissipate from (e) radiation pressure, which would prevent (e) more massive stars from forming, and consequently set (eb) an approximate upper limit for stellar mass black holes.

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| (2427) Observational evidence for stellar-mass black holes comes primarily from (e) X-ray astronomy. |
| (2428) A black hole in a close binary orbit can form (eb) an accretion disk (see Figure 1) of matter pulled off (er) a normal stellar companion. In its inspiral, disk material is heated by shearing, and becomes (=) a strong X-ray emitter. |
| (2429) If observations reveal a point-like X-ray emitter with (e&eb) a mass inferred from (e) orbital dynamics to be (=) above that possible for a neutron star, it becomes (=) a black hole candidate. |
| (2430) t present there are over a dozen such black hole candidates, including the first object to be identified as (=) a black hole, the X-ray source Cygnus X-1 (estimated mass 10±3Mo) [Casares (2006)]. |
| (2431) The other class of observationally-supported black holes is (=) that of supermassive black holes, ranging from hundreds to (e&eb) billions of solar masses. |
| (2432) Evidence for the existence of black holes at (eb) the upper end of this range is overwhelming. |
| (2433) By contrast, the existence of black holes with (e&*eb) mass roughly of order 102 to 104Mo, referred to as "intermediate-mass black holes," is speculative. |
| (2434) The supermassive holes may have begun as (=) primordial or stellar-mass black holes, but have grown through (e&eb) absorption of stars or gas or through mergers with (e&eb) other holes [Ferrarese and Ford (2005)]. |
| Figure 4: Observed orbits of stars near (eb) the dark compact object at the center of (e) our Milky Way galaxy: still frame from an animation. Animation created by Prof. Andrea Ghez and her research team at UCLA, from data sets obtained with the W. M. Keck Telescopes: http://www.astro.ucla.edu/~ghezgroup/gc/ |
| Evidence for supermassive black holes originally came from (e) quasars, small but intense radio sources seen at (eb) cosmological distances. |
| Their huge luminosities implied (eb) high mass, while (e) their rapid variability implied (eb) extremely small size. |
| More recently, measurements of Doppler shifts of stars and gas in the centers of galaxies have shown (eb) that compact objects of mass greater than 106Mo reside in (eb) the cores of most galaxies; the very small size of these central objects rules out any plausible alternative to the black hole explanation. |
| The black hole at the center of our own Milky Way galaxy has (e) a measured mass of ∼3×106Mo , based on (e) the orbits of stars near the radio source Sagittarius A* associated with (e&eb) the Galactic core [Ghez et al. (2005); see Fig. 4]. |
| (2440) Black holes are of particular interest for (e) gravitational waves, and vice versa. |
| (2441) For gravitationally-bound systems, the typical maximum gravitational wave amplitude (dimensionless strain) is (=) h∼(2GM/c2)2/(RD)∼R2H/(RD) , where D is the distance to the system, M and R are the mass and size of (e) the system, and RH is the horizon radius of (e) a black hole of the same mass. |
| (2442) From this it is clear that strongest gravitational waves will involve (e&eb) systems at or near black hole compactness: in particular, systems containing (er) black holes in close proximity to one another [Flanagan and Hughes (2005)]. |
| (2443) Conversely, a system containing only black holes would not be expected to radiate (eb) any sort of radiation other than gravitational waves, and in any system containing (e) black holes, gravitational waves provide (eb) the most accurate probes of conditions deep in (eb) their gravitational wells. |
| (2444) In particular, gravitational wave observations could definitively test (e&eb) whether a compact mass is truly a black hole [Cutler and Thorne (2002)]. |

(2445) The strongest intrinsic source of gravitational waves in the present Universe is (=) the collision of two comparable-massed black holes.

(2446) A sufficiently tight black hole binary system will lose (e) energy through emission of (e&eb) gravitational waves, causing (eb) the orbit to circularize and then slowly shrink over time.

(2447) Figure 5: Gravitational waves from (e) a black hole binary inspiral: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/

(2448) When the orbital radius approaches (eb) the Schwarzschild radius of the system, complex nonlinear dynamics come into play (e&eb) and the final stages of (e) the inspiral and merger can only be modeled through (e&eb) numerical simulations.

(2449) These simulations are quite challenging, but recent breakthroughs [Pretorius (2005)(See Fig. 5), Campanelli et al.(2006), Baker et al.(2006)] have led to (eb) the first complete inspiral-merger gravitational waveforms.

(2450) After the black holes have coalesced, they form (eb) a single highly-distorted black hole which quickly settles down to (e) a quiescent Kerr state through (e&eb) mission of quasinormal "ringdown" gravitational waves [Pretorius (2005), Campanelli et al. (2006), Baker et al. (2006)].

(2451) Another gravitational emission mechanism for supermassive black holes would come from (e) their capture of compact stellar objects, particularly stellar-mass black holes.

(2452) Unlike comparable-mass mergers, these captures would be on (eb) highly eccentric orbits that would not circularize before merger, resulting in (eb) intricate rapidly precessing orbits in (eb) the final months or years of inspiral.

**(2453)** The consequent gravitational-wave signal would allow (eb) for exquisitely precise measurements of the parameters of (e) the supermassive black hole, as well as tests of any deviation from (e) the predictions of general relativity [Cutler and Thorne (2002)].

**(2454)** For stellar-mass black holes (tens of Mo), the final stages of inspiral and merger occur at (eb) frequencies of hundreds to thousands of Hz, a regime targeted by (e) ground-based gravitational-wave observatories.

**(2455)** Systems of supermassive black holes that are (=) millions of Mo emit at frequencies of tens of (e) mHz, which will be detectable by (e) proposed space-based (e) gravitational-wave detectors.

**(2456)** Black holes in the billions of Mo, or a stochastic background of signals from (e) more moderate supermassive black holes at (eb) earlier stages of their inspiral, might emit (eb) waves at nHz frequencies that would be measurable through (e&eb) years-long correlated timing measurements of radio pulsars [Jenet et al. (2005)].

**An equation to quantify the origins of life on other planets July 5, 2016 by Bob Yirka report**

(2457) (Phys.org)—A pair of researchers, one with the Columbia Astrobiology Center in New York, the other with the University of Glasgow in the U.K. has come up with a mathematical equation that when (e) solved is meant to offer (eb) a means for estimating how often life begins on other planets. In their paper published in Proceedings of the National Academy of Sciences, Caleb Scharf and Leroy Cronin describe their equation, how they came up with it and why they believe it might become more useful as scientists learn more about the true nature of other planets and solar systems.

(2458) Back in the 1960's, an astronomer named Frank Drake came up with a formula for (e) estimating the likely number of (e) alien civilizations that might be capable of transmitting radio signals in such a way as to be (=) recognizable by receivers here on earth.

(2459) That equation, named quite naturally the **Drake Equation** has been the unofficial standard-bearer for half a century, despite it having no solution because (e) some of the parameters are still unknown.

(2460) The same holds true for the new equation developed by (e) Scharf and Cronin, but the results given should the unknowns ever be discovered would offer a different sort of answer than Drake was looking for.

(2461) To build their formula, the duo came up with several parameters to solve for (e) **abiogenesis**, which is the likelihood of an event occurring that leads to (eb) life beginning, they included (e) the number of possible building blocks, the mean number of such blocks per possible organism, the availability of building blocks that might exist during a given time period, expressed as a fraction and the probability that the existence of the building blocks would actually lead to (eb) life starting in a given unit of time. It looks like this: **Nabiogenesis (t)= Nb · 1/no · fc · Pa · t**

(2462) In essence the formula is meant to suggest (eb) that the probability of life beginning on a given planet is very likely connected to (e&eb) whether there are building blocks available on a given planet, and how much of them there might be.

(2463) As more is learned by space researchers it is hoped that the formula could help narrow down (e) search targets by offering (eb) a statistical probability of success for (e) a given planet.

(2464) The researchers note that the building blocks available do not necessarily have to be (=) of the type that led to (eb) life beginning on our planet.

**(2465)** They also suggest that implied in the formula is (=) the likelihood that life beginning events are (=) more likely to occur in solar systems where (e) there are multiple planets, which allow (eb) for an opportunity to share materials that could lead to (eb)building blocks **More information: Quantifying the origins of life on a planetary scale, Proceedings of the National Academy of Sciences, www.pnas.org/cgi/doi/10.1073/pnas.1523233113**

**Abstract**

(2466) A simple, heuristic formula with (e&eb) parallels to the Drake Equation is introduced to help focus discussion on open questions for (e) the origins of life in a planetary context.

(2467) This approach indicates (eb) a number of areas where (e) quantitative progress can be made on parameter estimation for (e) determining origins of life probabilities, based on (e) constraints from Bayesian approaches.

(2468) Authors discuss a variety of "microscale" factors and their role in determining (eb) "macroscale" abiogenesis probabilities on (e&eb) suitable planets.

(2469) Authors also propose that impact ejecta exchange between (e&eb) planets with parallel chemistries and (e&eb) chemical evolution could in principle amplify (eb) the development of (e) molecular complexity and (e&eb) abiogenesis probabilities.

(2470) This amplification could be very significant, and both bias conclusions about abiogenesis probabilities based on (e) the Earth and provide (eb) a major source of variance in the probability of life arising in (eb) planetary systems.

**(2471)** Authors use heuristic formula to suggest (eb) a number of observational routes for improving constraints on (e) origins of life probabilities. **Read more at: http://phys.org/news/2016-07-equation-quantify-life-planets.html#jCp**

**Brian Koberlein Shared publicly - Yesterday 8:00 PM the Illusion of Truth**

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| (2472) | Physical theories are often presented as (=) a description of what's really going on. |
| (2473) | Forces act on (e&eb) a baseball, causing (eb) baseball to fall. |
| (2474) | Atoms collide and fuse in (eb) the heart of a star, releasing heat and energy. Science is true, as is often said, and our scientific theories encapsulate this truth. But this isn't entirely true. |
| (2475) | For example, quantum mechanics is a strange, sometimes confusing theory. Objects can be (=) particles and waves. |
| (2476) | They do strange things like tunnel through (e&eb) barriers, and can even appear and disappear (e&eb). |
| (2477) | Given all this strangeness, what's really going on? What's the true nature of quantum reality? It depends on (=) which approach you want to take. |
| (2478) | In one view quantum objects are described by (e) a wave-like probability function. |
| (2479) | When you interact with (e&eb) or "observe" (e) the object this "wavefunction" collapses into (eb) a definite state. |
| (2480) | That's view commonly presented, but a quantum system can also be described by (e) its transitions between energy states. |
| (2481) | Since this uses a mathematical method involving (e&eb) matrices, it's known as (=) matrix mechanics. |
| (2482) | Both wavefunctions and matrix mechanics give (eb) the same results, but their view of what really goes on is very different. |
| (2483) | Then there's the path integral method |
| (2484) | . Rather than a wavefunction or matrix transitions, path integrals imagine quantum objects can take almost any path between (e&eb) two states. |
| (2485) | By summing all the possible paths you can derive (eb) the odds that it will occur. |
| (2486) | So which one is true? Are quantum objects distributed waves of (e) probability? |
| (2487) | Are they simply transitions between (e&eb) energy states? |
| (2488) | Do they take an infinite number of (e) paths between states? |
| (2489) | Each of these models makes (eb) the same predictions, so one version is no more "true" than the others. |
| (2490) | What happens in practice is that we'll use whatever method is useful at the time. They are equivalent models, so the best model for the job is the one we'll use. The only reason the wavefunction view is so common is that it's (=) the version usually taught to introductory students. You might think this uncertainty of truth is due to (e) the behavior of quantum physics itself. It's so strange and counterintuitive that we can't (e) wrap our puny brains around what's really going on. But the same thing occurs in lots of other fields. Even something as straight forward as basic Newtonian physics |
| (2491) | Toss a baseball in the air and (eb) Earth's gravitational force pulls it down. |
| (2492) | The force of gravity is a simple truth, right? While we often describe classical motion in terms of (e&eb) forces and acceleration, we can also describe it in terms of (e&eb) energy and momentum. |
| (2493) | In the Lagrangian and Hamiltonian approach, the path of a baseball is (=) the optimized path among possibilities. |
| (2494) | In this view a baseball's path is (=) the extrema of an energy equation, and force can be derived as a necessary consequence of (e) this. |
| (2495) | In the relativistic view the baseball follows (e&eb) a geodesic, which is (=) the minimal path through (e&eb) space and time. |

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| (2496) So is a baseball's motion due to (e) a gravitational force, an energy extrema, or a spacetime geodesic? |
| (2497) As with quantum theory, different approaches yield (eb) the same result. They are mathematically equivalent, so we can use whatever method is most useful at the time. |
| (2498) At its core, science is less about truth and more about models. The metaphysics underlying a model is useful only as far as it allows us (eb) to make better predictions, generate (eb) new ideas, or bring models together as (=) a cohesive whole |
| (2499) This is why we have no problem using (e) classical gravity to calculate the path of (e) a spacecraft through (e&eb) the solar system, while (e) using special relativity to account for (e) the Doppler shift of the spacecraft's radio signals. |
| (2500) It's why we can use (e) quantum physics to study atoms in the morning, and general relativity to (e) study black holes in the afternoon. |
| (2501) In regimes where models conflict with each other it isn't a failure of truth, but instead shows an opportunity to develop a better model. |
| (2502) It could be that with each better model we move closer to (e&eb) the truth about reality. |
| (2503) The truth is out there, and science strives to move towards (e&eb) that truth. |
| (2504) It's a common view, and certainly the search for truth has driven (eb) many scientists to develop better and better models. |
| (2505) But the real power of science is the recognition that what we have are models. Our models can be powerful, but they are always a bit tentative. There's always a chance that they might just be an illusion of truth. |

**Synopsis: A Single-Photon Cheshire Cat July 1, 2016 Researchers detected the polarization of a photon separate from the photon itself, just as the grin of Lewis Carroll's Cheshire cat can appear apart from the cat's body Synopsis figure APS/Carin Cain**

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| (2506) In a quantum Cheshire cat, one property of a particle is present detached from (e) the particle itself, just as (=) the grin of the feline in Alice in Wonderland can exist in the absence of (e) the whole animal. |
| (2507) Now a research team has successfully separated (e&eb) a single photon from (e) its polarization. |
| (2508) The effect was previously demonstrated (eb) with neutrons and also with light in a classical scenario, but the latest experiment uses single photons. |
| (2509) To produce the Cheshire cat, **Maximilian Schlosshauer and his colleagues** at the University of Portland in Oregon used an interferometer where (e) light hits a beam splitter, goes in two, opposite directions around a circuit (but along slightly offset routes), and then recombines at (eb) the entrance. |
| (2510) The incoming polarization of light was horizontal, but this was switched to (e&eb) vertical for the counterclockwise loop. |
| (2511) Each path included a device to adjust (e&eb) the polarization further and another to (e&eb) scatter away a small number of photons of (e) a designated polarization, the equivalent of detecting their presence. |
| (2512) The team fired single photons into (e&eb) the interferometer so that only one photon was present at (eb) a time, traversing both paths simultaneously. |
| (2513) They placed a horizontal polarizer in front of the output photon detector and gradually varied (e&eb) the difference between the two path lengths. |

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| (2514) This variation generated (eb) a sine-wave interference pattern from (e) the exiting photons only when the polarization was tweaked in (eb) the counterclockwise path. |
| (2515) At the same time, the average number of photons counted decreased only when (e) light was scattered from the clockwise path. |
| (2516) Simultaneously demonstrating these different effects for two routes of a single photon constitutes (eb) evidence of the Cheshire cat. |
| **(2517)** The effect might someday improve precision experiments if (e) an unwanted property could be moved away from the property being measured. **This research is published in Physical Review A. – David Ehrenstein David Ehrenstein is the Focus Editor for Physics.** |

**Observation of the quantum paradox of separation of a single photon from one of its properties James M. Ashby, Peter D. Schwarz, and Maximilian Schlosshauer Phys. Rev. A 94, 012102 – Published 1 July 2016**

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| (2518) Authors report an experimental realization of the quantum paradox of (e) the separation of a single photon from (e) one of its properties (the so-called "quantum Cheshire cat"). |
| (2519) Authors use a modified Sagnac interferometer with (e&eb) displaced paths to produce (eb) appropriately pre- and postselected states of (e) heralded single photons. |
| (2520) Weak measurements of photon presence and (e&eb) circular polarization are performed in each arm of (e) the interferometer by introducing (eb) weak absorbers and small polarization rotations and analyzing changes in (eb) the postselected signal. |
| (2521) The absorber is found to have (e) an appreciable effect only in one arm of the interferometer, while (e) the polarization rotation significantly affects (e&eb) the signal only when (e) performed in the other arm. |
| (2522) Authors carry out both sequential and simultaneous weak measurements and find (eb) good agreement between measured and predicted weak values. |
| (2523) In the language of Aharonov et al. and in the sense of the ensemble averages described by (e) weak values, the experiment establishes (eb) the separation of a particle from one its properties during the passage through (e&eb) the interferometer. |

**Viewpoint: Relativity Gets Thorough Vetting from LIGO Frans Pretorius, Department of Physics, Princeton University, Princeton, NJ 08544, USA May 31, 2016• Physics 9, 52 The gravitational-wave signal observed by the LIGO detectors shows no deviation from what general relativity predicts. Figure caption expand figure Frans Pretorius, APS/Carin Cain**

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| (2524) Figure 1: The signal from one of the LIGO detectors in Hanford, Washington, is shown with two representations of (e) the best-fit numerical relativity (NR) waveform. |
| (2525) The filtered NR waveform illustrates how the raw waveform is perceived by (e) the detector, |
| (2526) On September 14, 2015, the LIGO interferometers made the first direct measurement of (e) gravitational waves [1]. |
| (2527) The historic discovery came nearly 100 years after Einstein showed that his general relativity theory predicts (eb) the existence of gravitational waves and over 50 years since (e) the first quest to observe the Universe in gravitational waves [2]. |
| (2528) The observed event, labeled GW150914, was (=) the collision of two black holes, one weighing 29 solar masses, the other 36 solar masses. |

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| (2529) This binary coalesced into (e&eb) a remnant black hole weighing 62 solar masses and rotating (e&eb) at 67% the maximum rate allowed (eb) for a black hole. |
| (2530) The most statistically significant portion of (e) the signal came from (e) the last several gravitational-wave cycles of the collision, where (e) the most nonlinear, strong-field dynamics of the theory manifest. |
| (2531) One couldn't have asked for a better first event to learn about (e&eb) this regime of general relativity. |
| (2532) The first thing that GW150914 can teach us is how closely nature follows (eb) Einstein's prescriptions for (e&eb) colliding black holes. |
| (2533) Researchers from the LIGO team and the Virgo team have performed a series of tests to see if (e) any violations of (e) general relativity might be lurking in the data of GW150914 [3]. |
| (2534) They first determined (eb) the theoretical waveform, or "template," from (e) general relativity simulations that best fit the measured signal (see Fig. 1). |
| (2535) Their main finding is that the residual signal—obtained by (e) subtracting this template from (e) the data—is consistent with (e&eb) noise. |
| (2536) This is, in a sense, the "everything" test, as (e) the template folds in (e&eb) all that general relativity predicts about the merger: the decay of the last few orbits of (e) the two black holes, their violent collision (e&eb) and astonishingly rapid relaxation to (e) a quiescent rotating black hole, the propagation of (e) gravitational waves through (e&eb) the cosmos, and finally the transformation of (e) these waves into(e&eb) a chirp in the LIGO detectors. |
| (2537) Given the large signal-to-noise ratio of 24, the close agreement between the template and (=) the data puts limits on (e) alternative models. |
| (2538) Specifically, the predictions of these models can only disagree with (e&eb) the general relativity prediction for GW150914 by, at most, 4%. |
| (2539) In addition to this comprehensive test, the LIGO/Virgo teams also performed other checks that pick apart the signal, in order to address particular aspects of general relativity and black hole physics. In one such test, they compared the early and (e&eb) late parts of the waveform. |
| (2540) From the waves generated during (e&eb) the initial inspiral, one can infer (eb) the masses and spins of the binary. |
| (2541) Given these numbers, general relativity allows (eb) the mass and spin of (e) the remnant black hole to be calculated. |
| (2542) The researchers compared this result to (e&eb) an independent measurement of (e) the mass and spin obtained from (e) the late-time ringdown, and to within (eb) uncertainties, the two estimates were (=) consistent. |
| (2543) Another test isolated coefficients in the so-called Post Newtonian (PN) series describing (eb) the inspiral. |
| (2544) The PN series is (=) an expansion of the gravitational-wave signal in powers of (e) the velocity vv of the sources relative to (e&eb) the speed of light cc. |
| (2545) Observations of binary pulsars, which show (eb) a decrease in their orbital period due to (e) gravitational-wave emission, have constrained (e) PN terms, especially at low order. |
| (2546) For higher-order terms associated with (e&eb) strong-field emission, the LIGO/Virgo teams were able to improve (eb) the limits by many orders of magnitude compared to (e&eb) the binary pulsar constraints. |

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| (2547) At first glance it might seem these new results simply add to (e) the litany of previous tests that general relativity passed with flying colors. |
| (2548) However, GW150914 is vastly more significant than that. Other tests have shown that general relativity is valid in (eb)the weak-field regime, but GW150914 gives (eb) the first direct observational evidence that general relativity is the correct theory of nature governing (e&eb) the extreme spacetime dynamics around the event horizon of (e) a black hole. |
| (2549) One could argue that binary pulsar observations already told us that gravitational waves exist [4], but GW150914 is more luminous in gravitational waves than pulsar binaries by (e) at least 24 orders of magnitude. |
| (2550) It's true that the evolution of the Universe as a whole is driven by (e) strong gravity, but it would be difficult to argue that precision cosmological measurements can constitute (eb) a test of general relativity when (e) one of the dominant inputs, dark energy, is still a mystery. |
| (2551) As far as black holes are concerned, astronomers have (e) observed ultra compact dark objects that are consistent with (e&eb) black holes. |
| (2552) The most striking data comes from (e) stellar orbits about the 4 million solar mass objects in (eb) the center of our Galaxy [5], but present observations cannot constrain it to be (=) a black hole as described by general relativity. |
| (2553) GW150914 has given us the first morsels of quantifiable evidence that the rich physics of black holes is as predicted by general relativity. For example, the LIGO observations are consistent with (e&eb) so-called cosmic censorship [6], which says that gravitational singularities must remain hidden behind an (eb) event horizon. |
| (2554) If a "naked" singularity were produced in (eb) the energetic collision of GW150914, the data would likely have shown (eb) some sort of anomaly, but in exactly what form, we don't know. |
| (2555) Another example involves the "no-hair" theorem [7], which says that the only properties of a black hole are (=) its mass, spin, and charge (though astrophysical black holes are expected to have negligible charge). |
| (2556) By consequence, these properties should determine (eb) the frequency and decay constants of (e) all the infinitely many ringdown modes of (e) a black hole. |
| (2557) With GW150914, only the least-damped ringdown mode is (=) visible above the noise, but when combined with (e&eb) the remnant mass and spin inferred from (e) the inspiral, this is enough to count as (=) the first piece of evidence for a black hole's baldness. |
| (2558) GW150914 should also be a test for (e) alternative theories of gravity or exotic alternatives to (e) black holes within general relativity, such as boson stars [8]. |
| (2559) Some of these alternative models are concerned with (e&eb) explaining dark energy or dark matter and do not (e) therefore, predict any deviation from relativity at (eb) the scale of GW150914. |
| (2560) However, other models do in fact deviate from (e) conventional relativity in their predictions of (e) the gravitational-wave signal from (e) a binary merger. |
| (2561) The largest deviations should presumably occur during (e&eb) the last several gravitational-wave cycles, where (e) the signal is strong, and deviations should be (=) easy to detect or constrain. |
| (2562) Embarrassingly though, the majority of alternatives that claim deviations can't (e) really say much about this event, as the two-body problem hasn't been solved in these theories. |
| (2563) These models can only reproduce (eb) a fraction of the full gravitational-wave signal, (eb) so the strong constraints one might have hoped for (e) cannot be leveled. |

(2564) Exotic alternatives within (eb) general relativity are in a similar boat, or worse for (e) objects like gravastars, which have (e) no well-defined theoretical framework to say what happens when (e) two such objects collide.

(2565) Regardless, what is clear from the data is that any non-black-hole solution must imagine a form of (e) exotic matter with a viscosity larger than (e) any known material to explain (eb) the rapid damping of oscillations observed during the ringdown.

(2566) The inferred event rate for heavy black hole mergers like GW150194 is high, and LIGO has another factor of 3 improvement in sensitivity slated before it reaches (eb) design sensitivity.

(2567) Within 5 years then, we can expect dozens of detected events, with the loudest having a signal-to-noise ratio upwards of (e) 100.

(2568) This will allow (eb) for exquisite insight into the nature of extreme gravity beyond (e) what GW150194 has already gifted us with. **This research is published in Physical Review Letters**.

**Tests of General Relativity with GW150914 B. P. Abbott et al. (LIGO Scientific and Virgo Collaborations) Phys. Rev. Lett 116, 221101 – Published 31 May 2016**

(2569) The LIGO detection of GW150914 provides (eb) an unprecedented opportunity to study the two-body motion of (e) a compact-object binary in (eb) the large-velocity, highly nonlinear regime, and to witness the final merger of (e) the binary and the excitation of uniquely relativistic modes of (e) the gravitational field.

(2570) Authors carry out several investigations to determine whether GW150914 is consistent with (e&eb) a binary black-hole merger in general relativity.

(2571) They find that the final remnant's mass and spin, as determined from (e) the low-frequency (inspiral) and (e&eb) high-frequency (postinspiral) phases of the signal, are mutually consistent with (e&eb) the binary black-hole solution in general relativity.

(2572) Furthermore, the data following the peak of GW150914 are consistent with (e&eb) the least-damped quasinormal mode inferred from (e) the mass and spin of the remnant black hole.

(2573) By using waveform models that allow (eb) for parametrized general-relativity violations during (e&eb) the inspiral and merger phases, authors perform quantitative tests on (e&eb) the gravitational-wave phase in the dynamical regime and they determine the first empirical bounds on (e&eb) several high-order post-Newtonian coefficients.

(2574) Authors constrain (e) the graviton Compton wavelength, assuming (eb) that gravitons are dispersed in (eb) vacuum in the same way as particles with mass, obtaining a 90%-confidence lower bound of 1013 km.

(2575) In conclusion, within our statistical uncertainties, we find no evidence for violations of (e) general relativity in the genuinely strong-field regime of gravity. **Received 26 March 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.221101 © 2016 American Physical Society**

**Focus: Background Noise of Gravitational Waves March 31, 2016• Physics 9, 33**

(2576) Assuming LIGO's recently detected black hole merger is not unusual; researchers revised upward their estimate of (e) the strength of the background noise coming from (e) distant mergers across the Universe. Figure caption expand figure Caltech/MIT/LIGO Lab (deleted due to spatial constraints)

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| (2577) Catching some waves The LIGO-Livingston interferometer consists of (e) two 4-km-long arms where (e) laser light bounces back and forth. Interference between the lasers from the two arms can signal tiny changes in (e&eb) the lengths of the arms resulting from gravi... |
| (2578) The recent detection of gravitational waves emitted from (e) a pair of merging black holes implies (eb) the existence of a large population of similar pairs, which should be generating (eb) a background of waves too faint to be detected with (e&eb) present technology. |
| (2579) Using this information, researchers now estimate that this background could be (=) roughly 10 times stronger than previously expected and might be found by (e) sensitive instruments in only a few years. |
| (2580) The background signal could help astrophysicists better understand how black hole binaries form and also pave the way to detecting the waves emitted by (e) the expanding Universe shortly after the big bang. |
| (2581) Gravitational waves were detected for the first time on 14 September 2015 by the Laser Interferometer Gravitational-Wave Observatory (LIGO). |
| (2582) By reflecting laser light between mirrors set 4 km apart, the team was able to detect (e&eb) tiny fluctuations in the distance between the mirrors as (=) the waves passed. |
| (2583) The waves were triggered by (e) an event, called GW150914; involving (e&eb) the merger of a pair of black holes having (e) masses of about 29 and 36 times that of the Sun. |
| (2584) Researchers expect that similar pairs of black holes, as well as neutron stars, should be merging constantly, creating (eb) gravitational waves that should bombard (e&eb) the Earth from all directions and make the equivalent of (=) a faint popping noise in a detector. |
| (2585) Previous estimates of the likely strength of this background had to make assumptions about (e&eb) how many pairs there might be and about the distribution of their masses. |
| (2586) These estimates remained speculative, as no actual binaries were known. |
| (2587) "What the recent observation has changed," says Nelson Christensen of Carleton College in Minnesota and the LIGO Scientific Collaboration (LSC), "is that we now know there really is (=) a population of black hole binaries. We have a small bit of information about it." |
| (2588) The new analysis—carried out by LSC researchers working with (e&eb) the European Virgo collaboration—estimates (eb) the energy density in the gravitational wave background in (eb) the frequency range below 100 Hz, where (e) the LIGO detector is most sensitive. |
| (2589) To make this estimate, they assumed (eb) a population of black hole binaries of various sizes, having (e) a range of distances from Earth, and created through (e&eb) a variety of scenarios. |
| (2590) They then calculated the amount of gravitational radiation that would result from (e) the merging of these objects. |
| (2591) The team then adjusted this estimate by using (e) details of the GW150914 event—especially the masses of the black holes involved—to make (eb) a most likely guess at the actual distribution of black hole binaries. |
| (2592) The idea is that the observed merger is (=) most likely representative of (e) the general population, rather than being extremely unusual. |
| (2593) If the black hole binary population is represented as (=) a bell curve, or Gaussian, the event should lie within (eb) the main bulge, not in the tails. |
| (2594) The observed mass of GW150914 suggests (eb) that the main population of black hole binaries could be (=) roughly 20 times more massive than expected, which would generate (eb) a background roughly 10 times "louder" than (e) previous estimates. |

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| (2595) However, Christensen emphasizes that the result has (e) lots of uncertainty, as it's based on (e) only one event. |
| (2596) "This estimate seems very plausible to me," says astrophysicist Maura McLaughlin of the University of West Virginia in Morgantown. She is a member of the NANOGRAV collaboration, which tries to detect gravitational waves at much lower frequencies than LIGO can observe. "I can't see any holes in their arguments. Unless LIGO got really lucky, then there must be many more of these systems than previously thought." This level of gravitational background implies (eb) that the LIGO and Virgo detectors might be sensitive enough to see it in 2–5 years, sooner than physicists previously expected, says Christensen, which has led to much anticipation. "Everyone is super excited," he says. |
| (2597) Detecting the background would provide (eb) a confirmation of the presence of gravitational waves carrying (e&eb) energy throughout the Universe. It would also be useful, the researchers suggest, in gaining (e) further understanding of precisely how black holes form. |
| (2598) And subtracting this background might someday allow (eb) detection of the fainter waves coming from (e) the expansion of the Universe moments after the big bang. This research is published in Physical Review Letters. –Mark Buchanan Mark Buchanan is a freelance science writer in Normandy, France. |

**GW150914: Implications for the Stochastic Gravitational-Wave Background from Binary Black Holes B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 131102 – Published 31 March 2016**

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| (2599) The LIGO detection of the gravitational wave transient GW150914, from (e) the inspiral and merger of (e) two black holes with masses $\gtrsim 30 M_{\odot}$, suggests (eb) a population of binary black holes with (e&eb) relatively high mass. |
| (2600) This observation implies (eb) that the stochastic gravitational-wave background from binary black holes, created from (e) the incoherent superposition of (e) all the merging binaries in the Universe, could be (=) higher than previously expected. |
| (2601) Using the properties of GW150914, authors estimate the energy density of (e) such a background from (e) binary black holes. |
| (2602) In the most sensitive part of the Advanced LIGO and Advanced Virgo band for (e) stochastic backgrounds (near 25 Hz), authors predict $\Omega_{GW}(f=25\ \mathrm{Hz}) = 1.1^{+2.7}_{-0.9} \times 10^{-9}$ with 90% confidence. |
| (2603) This prediction is robustly demonstrated for (e) a variety of formation scenarios with (e&eb) different parameters. |
| (2604) The differences between models are (=) small compared to the statistical uncertainty arising from (e) the currently poorly constrained (e) local coalescence rate. |
| (2605) Authors conclude that this background is potentially measurable by (e) the Advanced LIGO and Advanced Virgo detectors operating at their projected final sensitivity. DOI:http://dx.doi.org/10.1103/PhysRevLett.116.131102 © 2016 American Physical Society |

**Study opens new realms of light-matter interaction July 15, 2016 by David L. Chandler. Some sentences might have been omitted due to spatial constraints. Please bear with me.**

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| (2606) Emission spectra are a widely used method for (e) identifying chemical compounds; the bright lines reveal (eb) the different frequencies of light that can be emitted by (e) an atom. |

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| (2607) Here, a normal emission spectrum for (e) an atom in a high-energy state (top) is compared to (e&eb) the emission from the same atom placed just a few nanometers (billionths of a meter) away from (e) graphene that has been doped with (e&eb) charge carriers (bottom). |
| (2608) For each energy-level transition, an orange line (or purple cloud) appears if (e) that transition is estimated to be faster than (e) one per microsecond — making it frequent enough to be (=) observed. **Credit: Massachusetts Institute of Technology** |
| (2609) A new MIT study could open up new areas of technology based on (e) types of light emission that had been thought to be (=) "forbidden," or at least so unlikely as to be (=) practically unattainable. |
| (2610) The new approach, the researchers say, could cause (eb) certain kinds of interactions between light and (e&eb) matter, which would normally take billions of years to happen, to take place instead (e&eb) within billionths of a second, under (e&eb) certain special conditions. The findings, based on a theoretical analysis, are reported today in the journal Science in a paper by MIT doctoral student Nicholas Rivera, Department of Physics Professor Marin Soljačić, Francis Wright Davis Professor of Physics John Joannopoulos, and postdocs Ido Kaminer and Bo Zhen. |
| (2611) Interactions between light and matter, described by (e) the laws of quantum electrodynamics, are the basis of a (e) wide range of technologies, including (e) lasers, LEDs, and atomic clocks. |
| (2612) But from a theoretical standpoint, "Most light-matter interaction processes are 'forbidden' by (e) electronic selection rules, which limits (e) the number of transitions between energy levels we have access to," Soljačić explains. |
| (2613) For example, spectrograms, which are used to (e) analyze the elemental composition of materials, show (eb) a few bright lines against (e&eb) a mostly dark background. |
| (2614) The bright lines represent (eb) the specific "allowed" energy level transitions in (eb) the atoms of that element that can be accompanied by (e) the release of a photon (a particle of light). |
| (2615) In the dark regions, which make up most of the spectrum, emission at those energy levels is (=) "forbidden." |
| (2616) With this new study, Kaminer says, "we demonstrate theoretically that these constraints can be lifted" using (e) confined waves within (e&eb) atomically thin, 2-D materials. |
| (2617) "We show that some of the transitions which normally take the age of the universe to happen could be made to happen within (eb) nanoseconds. |
| (2618) Because of this, many of the dark regions of a spectrogram become (=) bright once an atom is (=) placed near a 2-D material." |
| (2619) Electrons in an atom have (e) discrete energy levels, and when they hop from one level to another they give off (eb) a photon of light, a process called spontaneous emission. |
| (2620) But the atom itself is much smaller than the wavelength of the light that gets emitted—about 1/1,000 to 1/10,000 as big—substantially impairing (e&eb) the interactions between the two. |
| (2621) The trick is, in effect, to "shrink" (e) the light so it better matches the scale of the atom, as (=) the researchers show in their study. |
| (2622) The key to enabling a whole range of interactions, specifically transitions in (eb) atomic states that relate to absorbing or emitting light, is (=) the use of a two-dimensional material called graphene, in which light can interact with (e&eb) matter in the form of plasmons, a type of electromagnetic oscillation in the material. |
| (2623) These plasmons, which resemble photons but have wavelengths hundreds of times shorter, are (=) very narrowly confined in the graphene, in a way that makes some kinds of interactions with (e&eb) that matter many orders of magnitude more likely than (e) they would be in ordinary materials. |

(2624) This enables a variety of phenomena normally considered unattainable, such as (e) the simultaneous emission of multiple plasmons, or two-step light-emitting transitions between (e&eb) energy levels, the team says.

(2625) This method can enable the simultaneous emission of (e) two photons that are "entangled," meaning (eb) they share the same quantum state even when (e) separated.

(2626) Such generation of entangled photons is (=) an important element in quantum devices, such as those that might be used for cryptography.

(2627) Making use of these forbidden transitions could open up (eb) the ability to tailor the optical properties of materials in ways that had not been thought possible, Rivera says.

(2628) "By altering these rules" about the relationship between light and (e&eb) matter, "it can open (eb) new doors to reshaping (e&eb) the optical properties of materials."

(2629) Kaminer predicts that this work "will serve as (=) a founding piece for (e) the next generation of studies on light-matter interactions" and could lead to (eb) "further theoretical and experimental advances in many fields which rely on (e) light-matter interactions, including (e) atomic, molecular and optical physics, photonics, chemistry, optoelectronics, and many others."

(2630) Beyond its scientific implications, he says, "this study has possible applications across multiple disciplines, since (e) in principle it has potential to enable the full use of (e) the periodic table for optical applications."

(2631) This could potentially lead to (eb) applications in spectroscopy and sensing devices, ultrathin solar cells, and new kinds of materials to absorb solar energy, organic LEDs with higher efficiencies, and photon sources for (e) possible quantum computing devices."From the standpoint of fundamental science, this work lays the groundwork for a subfield that just a few years ago was difficult to imagine and until now was largely unexplored," Soljačić says.

(2632) "Two-dimensional materials confine (e) fields to a surface and motion to a plane, making possible (eb) many effects that are orders-of-magnitude too weak to (e) appear in a bulk volume," says Jason Fleischer, an associate professor of electrical engineering at Princeton University, who was not involved in this research.

(2633) This work, he says, "systematically explores how 2-D materials improve light-matter interactions, laying a theoretical foundation for faster electronic transitions, enhanced sensing, and better emission, including the compact generation of broadband and quantum light." More information: N. Rivera et al: Shrinking light to allow (eb) forbidden transitions on the atomic scale, Science (2016). DOI: 10.1126/science.aaf6308 Journal reference: Science search and more info website: Provided by: Massachusetts Institute of Technology search and more info website. This story is republished courtesy of MIT News (web.mit.edu/news office/), a popular site that covers news about MIT research, innovation and teaching.

**Viewpoint: Peering Back to the Beginning of Time Lawrence M. Krauss, Department of Physics and School of Earth and Space Exploration, Arizona State University, Tempe, AZ 85287, USA and Mount Stromlo Observatory, Research School of Astronomy and Astrophysics, Australian National University, Weston, ACT 2611, Australia June 19, 2014• Physics 7, 64**

(2634) The BICEP2 collaboration reports the detection of BB-mode polarization of (e) the cosmic microwave background—a signal that might originate from (e) gravitational waves created by (e) inflation during the very earliest moments in the evolution of our Universe. Figure caption expand figure P. A. R. Ade et al. [1]; Image on homepage: BICEP2 Collaboration (deleted due to restrictions on space)

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| (2635) Figure 1: (Top) BB-mode polarization pattern observed by (e) the BICEP2 telescope [1]. |
| (2636) The twisting pattern in the polarization of (e) the cosmic microwave background (CMB) might originate from (e) gravitational waves created by inflation in the very first moments of the universe. |
| (2637) Inflation—the hypothesis that the Universe underwent (e&eb) a phase of superluminal expansion in a brief period following the big bang—has (e) the potential of explaining, from first principles, why the Universe has (e) the structure we see today. |
| (2638) It could also solve outstanding puzzles of standard big-bang cosmology, such as why the Universe is, to a very good approximation, (=) flat and isotropic (i.e., it looks the same in all directions). |
| (2639) Yet we do not yet have a compelling model, based on (e) fundamental particle physics principles, that explains inflation. |
| (2640) And despite its explanatory power and a great deal of suggestive evidence, we still lack an unambiguous and direct probe of (e) inflation. |
| (2641) Theorists have developed different models for inflation, which all share (e&eb) a common, robust prediction: Inflation would have created (eb) a background of gravitational waves that could have (eb) an observable effect. |
| (2642) These waves would cause (eb) subtle, characteristic distortions of (e) the cosmic microwave background (CMB)—the oldest light in the Universe, released when (e) photons decoupled from matter and the Universe became (=) transparent to radiation |
| (2643) Until now, such a background of gravitational waves has (e) escaped detection. It is therefore very exciting that the BICEP2 (Background Imaging of Cosmic Extragalactic Polarization) collaboration has now reported in Physical Review Letters the observation of polarization modes of CMB radiation [1] that are consistent with (e&eb) such a background. |
| (2644) If confirmed, BICEP2's finding would provide (eb) the first direct probe of physical processes that occurred a mere 10−3610-36 seconds after (e) the big bang. |
| (2645) The result also offers a rare opportunity to directly test theoretical ideas, including inflation and grand unification (the theory that merges the electromagnetic, weak, and strong forces into a single force), at energies 13 orders of magnitude greater than (e) we can currently access at the Large Hadron Collider. |
| (2646) Inflation is based on (e) the possibility that a scalar field in the early Universe got momentarily stuck in (eb) a "false vacuum state"—a metastable condition in which the Universe wasn't at (eb) its global potential minimum. |
| (2647) The energy stored in this vacuum state could drive (eb) an exponential expansion of spacetime (similar to, but much faster than, the expansion due to dark energy seen today) before (e) a phase transition to the true vacuum state occurs, releasing (eb) this energy into (e&eb) particles and radiation. |
| (2648) This creates (eb) an initially hot state for (e) a subsequent big-bang-like expansion from (e) which our Universe subsequently evolved [2]. |
| (2649) If the Universe grew by at least a factor of 10261026, (eb) inflation can explain (eb) both the apparent flatness and isotropy of the Universe [2]: |
| (2650) The expansion would stretch (eb+ eb-dot) any initial curvature of (e) the Universe to near flatness, and smooth out (e&eb) any initial inhomogeneities, just as stretching (eb+ ebdot) a piece of fabric would do. |
| (2651) Isotropy is justified by the fact that the Universe, before inflation, was once much smaller and could have been contained in (eb) a causally connected region that was in thermal equilibrium, resulting in (eb) the approximately uniform, isotropic temperature distribution observed on (eb) large scales today. |

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| (2652) Importantly, the theory can explain the (eb) origin of structure in the Universe: Microscopic quantum fluctuations associated with (e&eb) the scalar field driving inflation could have been "frozen in" during the exponential expansion, resulting, (eb) after inflation, in density fluctuations that might have seeded all observed large-scale structures in the Universe, including (e) all visible galaxies and clusters of galaxies. |
| (2653) The signs of such primordial fluctuations have indeed been detected through (e&eb) their effect on (e&eb) the CMB, first by the Cosmic Background Explorer (COBE) [3] and later by other missions. |
| (2654) These observational studies detected small temperature variations (anisotropies) in (eb) the CMB and found a spectrum compatible with (e&eb) that predicted by inflation. |
| (2655) The magnitude of these fluctuations suggests (eb) an energy scale of inflation (loosely speaking, the energy stored throughout space in (eb) the scalar field potential that drove inflation) close to (e) 10161016 giga-electron-volts (GeV). |
| (2656) But while consistent with inflation, such fluctuations are not (e) conclusive evidence: Their predicted shape and magnitude is known to vary with (e&eb) the details of the inflationary potential (whose form is not yet experimentally constrained). |
| (2657) Hence many different observations could be accommodated, and any agreement between theory and observation is at best suggestive. There is also a possibility that other, noninflationary sources might contribute. However, gravitational waves—the propagating ripples in spacetime predicted by Einstein in 1916—might provide (eb) direct evidence of inflation. |
| (2658) During inflation, gravitational fields would also fluctuate because of (e) quantum effects. |
| (2659) These fluctuations would similarly get "frozen in" and re-emerge after inflation as (=) classical gravitational waves, whose amplitude depends only on (=) the energy stored in the field driving (eb) inflation, and not on other model-dependent details [4]. |
| (2660) If the scale of inflation were indeed as high as a few times 10161016 GeV, this background could produce (eb) specific temperature anisotropies in the CMB [5]. |
| (2661) These anisotropies have thus far not been detected by (e) the most recent searches by the WMAP and PLANCK satellites, and their absence was used to (e) put upper limits on the scale of inflation somewhat below 10161016 GeV. |
| (2662) But a more sensitive CMB probe of (e) gravitational waves exists. |
| (2663) Large-wavelength gravitational waves produced by (e) inflation would have alternatively compressed and expanded (e&eb) space, producing (eb) spatial variations in (eb) the intensity of radiation bombarding free electrons in (eb)the early Universe (see, e.g., Ref: [6]). |
| (2664) As photons decoupled from (e) matter and CMB radiation was released, scattering from (e) the undulating sea of electrons would have created (eb) patterns in the polarization of CMB. |
| (2665) Such polarization anisotropies are predicted in (eb) inflationary models to be (=) most visible when (e) looking at the sky on spatial scales comparable to (e&eb) or somewhat larger than the size of (e) the Universe's horizon at (eb) the time when the CMB was created. |
| (2666) This corresponds to (e&eb) about 1 degree of the field of view of the sky. |
| (2667) Over the past decade, various CMB polarization experiments have thus chased the observation of polarization anisotropies on (e&eb) these scales but, so far, also yielded (eb) merely upper limits on a possible gravitational background. |
| (2668) All of this has now changed with (e&eb) the BICEP2 collaboration's [1] report of (e) a signal that has the predicted (eb) characteristics of such a background at precisely the angular scales where (e) the inflationary signal is expected (see Fig. 1, top deleted due to spatial restraints). |

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| (2669) The BICEP2 Collaboration uses (e) a microwave telescope located at the South Pole, with (e&eb) an array of polarization-sensitive detectors pointing at (eb) a small field (a solid angle of about 11 percent of the full sky) with a resolution of approximately 0.50.5 degrees, thus sufficient to probe the relevant angular scales for the detection of (e) polarization anisotropies originating from (e) **inflation**. |
| (2670) BICEP2's search is based on (e) two expected features of primordial gravitational waves. |
| (2671) First, gravitational waves would produce (eb) a characteristic spatial polarization pattern on CMB—a consequence of (e) their quadrupole nature. |
| (2672) One can decompose (e&eb) any spatial polarization pattern into (e&eb) two orthogonal components, a "curl-free" component and (e&eb) a component involving a twisting pattern. |
| (2673) In analogy to electromagnetic modes, these polarization patterns are called "EE modes" and (e&eb) "BB modes," respectively (see Fig. 1). |
| (2674) Scalar fluctuations in the density of (e) matter would produce (eb) EE modes, while (e) gravitational waves would instead generate (eb) a BB-mode pattern. |
| (2675) The second feature is the aforementioned characteristic spatial scale of (e) expected CMB polarization signatures, including (e) wavelength scales comparable to (e&eb) or larger than the horizon scale at the time the CMB was created. |
| (2676) Since inflation would produce (eb) both gravitational waves and (e&eb) scalar density perturbations, BICEP2, like all CMB-polarization experiments, quote their results in terms of (e&eb) the ratio rr of the amplitude of (e) tensor perturbations (associated with (e&eb) gravitational waves) to (e) the amplitude associated with (e&eb) scalar density fluctuations (reflected in (eb) the magnitude of CMB temperature fluctuations measured by (e) experiments like COBE). |
| (2677) A nonzero rr implies (eb) the detection of a BB-mode pattern and primarily measures (eb) the strength of any possible gravitational wave signal which, if (e) it comes from inflation, is then determined by (e) its energy scale (e.g., see Ref. [7] for a review). |
| (2678) BICEP2 reported a nonzero value of r=0.2(+0.07−0.05)r=0.2(+0.07-0.05), on scales ranging from (e) the horizon scale up to 4–5 times larger, with (e&eb) a claimed significance of better than 5σ5σ. |
| (2679) If confirmed, a value of 0.2 corresponds to (e&eb) an energy density during (e&eb) inflation of ∼1.2×1016∼1.2×1016 GeV. |
| (2680) The claimed signal is far larger than expected, partly because (e) the PLANCK satellite, despite not having yet reported its polarization measurements, has (e) previously put an upper limit on (eb) rr (<0.11<0.11). PLANCK's estimate is based on (e) the absence of a gravitational wave signature on (eb) the CMB temperature anisotropy on (eb) its largest scales (PLANCK observes the whole sky and therefore angular scales up to 360360 degrees). |
| (2681) The apparent discrepancy between the PLANCK upper bound and (e&eb) the BICEP2 detection might be a cause for concern. |
| (2682) However, the two bounds probe different spatial scales, and the BICEP collaboration has noted that, while tensor perturbations are expected to be (=) roughly scale independent, some scale dependence of (e) scalar density perturbations could relax (e&eb) the PLANCK constraint and bring (eb) the two results into agreement. |
| (2683) Of greater concern, of course, is the possibility that the signal could be determined by (e) foreground contamination—any source along (e&eb) the line of sight of the detector that could distort or contaminate (e&eb) the primordial CMB signal. |
| (2684) BICEP2 tested for (e) a wide variety of possible foregrounds, claiming (eb) they cannot be responsible for (e) the observed signal. |

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| (2685) One obvious candidate arises from (e) gravitational lensing: The intervening mass of (e) astrophysical bodies could distort (e&eb) the spatial pattern of polarization in a way that that could transform (e&eb) an EE-mode signal into (e&eb) a BB-mode signal. |
| (2686) The BICEP collaboration has analyzed this possibility extensively. Based on (e) measurements of (e) CMB temperature (and hence density perturbations) and on (e&eb) models of the mass distribution, they argue that such effects would only be (=) significant on smaller angular scales than those contributing to their measurements of rr. |
| (2687) But since BICEP2's first announcement on March 17, 2014, a number of scientists have raised another foreground concern due to (e) the possible impact of (e&eb) our Galaxy's dust. |
| (2688) BICEP2 modeled dust contamination scenarios based on (e) a variety of theoretical models, using data from (e) its predecessor BICEP1 and preliminary dust estimates from (e) the PLANCK satellite. |
| (2689) They concluded that dust could not (e) reproduce the magnitude of the observed signal. |
| (2690) However, PLANCK has in the meanwhile released (eb) more complete measurements of variation in polarized dust foregrounds in (eb) many regions of the Galaxy. |
| (2691) Their analysis increases the uncertainty associated with (e&eb) the dust contribution to (e) the BICEP2 result and also accommodates (e) a potentially larger polarized dust contribution. |
| (2692) Some researchers argued that dust polarization at a level 2–3 times higher than that assumed by BICEP2 is possible and could therefore plausibly be responsible for (e) the entire signal (e.g., see Ref: [8]). |
| (2693) While the BICEP2 collaboration continues to maintain that this is extremely unlikely based on (e) their estimates, many other researchers are more skeptical. |
| (2694) In any case, a variety of upcoming measurements, from PLANCK and other detectors are (=) expected to resolve this important ambiguity within the coming year. |
| (2695) Even if the BICEP2 signal does arise from (e) primordial gravitational waves, various consistency checks will be required (e) before it can be definitively argued that the signal arises from (e) inflation. |
| (2696) Inflation is certainly the best-motivated possibility, and the spectral shape of the signal, at least for (e) the five largest-scale modes observed by BICEP2, agrees extremely well with (=) inflationary predictions (see Fig. 1, bottom). |
| (2697) Several results have already appeared that strongly constrain (e) other possible sources, ruling out (e) for example, the possibility of gravitational waves generated by (e) other possible phase transitions after inflation [9] |
| (2698) . This conclusion is reinforced by more careful analyses of the scale dependence of (e) the predicted signal (e.g., Ref. [10]). |
| (2699) But it will nevertheless be important to assure that all possible alternative sources of (e) gravitational waves are ruled out. |
| (2700) If it can be established that the BICEP2 represents (eb) a gravitational wave signal from (e) inflation, there is one additional exciting implication. |
| (2701) Standard calculations suggest (eb) that the gravitational wave signal from inflation is due to (e) quantum fluctuations in the gravitational field, and this appears to require (e) the existence of gravitons—the hypothetical elementary particles that mediate (e&eb) the force of gravity. |
| (2702) Research by Frank Wilczek and myself, based on (e) a dimensional analysis argument [11], suggests that,(eb) independently of the method used to calculate the gravitational wave spectrum, a quantum gravitational origin is (=) required. |
| (2703) If so, the BICEP2 result then implies (eb) that gravity is ultimately a quantum theory, a result of fundamental importance for physics. |

(2704) It is hard to imagine a more important or wider-ranging cosmological observation. If confirmed by upcoming experiments, the BICEP2 results will push the frontiers of empirical physics forward in an unprecedented way, offering (eb) the first direct observation of gravitational waves and of quantum gravity effects as well as a new probe of the earliest moments in the Universe's life and of physics at grand-unified energy scales. This will thus constrain both our understanding of fundamental interactions and the evolution of our Universe as a whole.

**Detection of B-Mode Polarization at Degree Angular Scales by BICEP2 P. A. R. Ade et al. (BICEP2 Collaboration) Phys. Rev. Lett 112, 241101 – Published 19 June 2014**

(2705) Authors report results from the BICEP2 experiment, a cosmic microwave background (CMB) polarimeter specifically designed to search for (e) the signal of inflationary gravitational waves in (eb) the B-mode power spectrum around $\ell \sim 80$.

(2706) The telescope comprised a 26 cm aperture all-cold refracting optical system equipped with (e&eb) a focal plane of 512 antenna coupled transition edge sensor 150 GHz bolometers each with temperature sensitivity of $\approx 300\ \mu K_{CMB}\sqrt{s}$. BICEP2 observed from (e) the South Pole for three seasons from 2010 to 2012.

(2707) A low-foreground region of sky with (e&eb) an effective area of 380 square deg was observed to (e) a depth of 87 nK deg in Stokes Q and U.

(2708) In this paper authors describe the observations, data reduction, maps, simulations, and results. They find an excess of B-mode power over the base lensed-ΛCDM expectation in the range $30<\ell<150$, inconsistent with (e&eb) the null hypothesis at (eb) a significance of $>5\sigma$.

(2709) Through jackknife tests and simulations based on (e) detailed calibration measurements authors show that systematic contamination is (=) much smaller than the observed excess.

(2710) Cross correlating against WMAP 23 GHz maps we find that Galactic synchrotron makes a negligible contribution to the observed signal. We also examine a number of available models of polarized dust emission and find that at their default parameter values they predict power ~(5–10)× smaller than the observed excess signal (with no significant cross-correlation with our maps). However, these models are not sufficiently constrained by external public data to exclude the possibility of dust emission bright enough to explain the entire excess signal. Cross correlating BICEP2 against 100 GHz maps from the BICEP1 experiment, the excess signal is confirmed with $3\sigma$ significance and its spectral index is found to be consistent with that of the CMB, disfavoring dust at $1.7\sigma$. The observed B-mode power spectrum is well fit by a lensed-ΛCDM+tensor theoretical model with tensor-to-scalar ratio $r=0.20+0.07-0.05$, with $r=0$ disfavored at $7.0\sigma$. Accounting for the contribution of foreground, dust will shift this value downward by an amount which will be better constrained with upcoming data sets. Received 4 April 2014 DOI:http://dx.doi.org/10.1103/PhysRevLett.112.241101 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**(2711) Physicist claims to have observed quantum effects of Hawking radiation in the lab for the first time August 16, 2016 by Bob Yirka report**

(2712) This computer-simulated image shows a supermassive black hole at the core of a galaxy. The black region in the center represents the black hole's event horizon, where no light can escape the massive object's gravitational grip. The black …more

(2713) (Phys.org)—Jeff Steinhauer, a physicist at the Israel Institute of Technology, has published a paper in the journal Nature Physics describing experiments in which he attempted to create a virtual black hole in the lab in order to prove that Stephen Hawking's theory of radiation emanating from black holes is correct —though his experiments are based on sound, rather than light. In his paper, he claims to have observed the quantum effects of Hawking radiation in his lab as part of a virtual black hole—which, if proven to be true, will be the first time it has ever been achieved.

(2714) For many years, scientists believed that nothing could ever escape from a black hole. But in 1974, Stephen Hawking published a paper suggesting that something could—particles that are now called Hawking radiation. His idea was that if a particle (and its antimatter mate) appeared spontaneously at the edge of a black hole, one of the pair might be pulled into the black hole while the other escaped, taking some of the energy from the black hole with it—which would explain why black holes grow smaller and eventually disappear. Because such emissions are so feeble, no one has been able to measure Hawking radiation, so researchers have instead tried to build virtual black holes in labs to test the theory. One type of virtual black hole was proposed back in 1981 by Bill Unruh with the University of British Columbia—he suggested that an analogue might be created using water instead of light. He imagined a phonon existing at the edge of a waterfall—as the water speeds up, it begins to move faster than the speed of sound, causing it to be trapped. But if the phonon had an entangled mate that eluded the fall by moving away before getting caught up, it could escape. In this new effort, Steinhauer has built a device based on that idea and in so doing, claims he has observed an analogue of Hawking radiation.

(2715) The experiment consisted of creating an entangled pair of phonons sitting inside a bit of liquid that had been forced (via laser) to move very fast and then observing the action as one of the pair was pulled away as the liquid began to move faster than the speed of sound, while the other escaped—the fluid was a Bose-Einstein condensate of rubidium-87 atoms. After repeating the experiment 4,600 times Steinhauer became convinced that the particles were entangled, a necessity for a Hawking radiation analogue. His findings do not prove Hawking's theory to be true, of course, but they do appear to add a degree of credence that other researchers have thus far not been able to achieve. More information: Jeff Steinhauer. Observation of quantum Hawking radiation and its entanglement in an analogue black hole, Nature Physics (2016). DOI: 10.1038/nphys3863

**(2716) Abstract**

We observe spontaneous Hawking radiation, stimulated by quantum vacuum fluctuations, emanating from an analogue black hole in an atomic Bose–Einstein condensate. Correlations are observed between the Hawking particles outside the black hole and the partner particles inside. These correlations indicate an approximately thermal distribution of Hawking radiation. We find that the high-energy pairs are entangled, while the low-energy pairs are not, within the reasonable assumption that excitations with different frequencies are not correlated. The entanglement verifies the quantum nature of the Hawking radiation. The results are consistent with a driven oscillation experiment and a numerical simulation. Journal reference: Nature Physics search and more info website

Read more at: http://phys.org/news/2016-08-physicist-quantum-effects-hawking-lab.html#jCp

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as

the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t - T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

These equations can be written for

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(19) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.

(20) However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].

(21) Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.

Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

**TIME DELAY**

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period,

which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types system takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t - T_7)F_7(R_7(t - T_7) - N_7(t)M_7((R_7(t) \qquad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7}+\frac{N_7 dG_7}{dR_7}-\frac{dB_7}{dR_7}>\frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7}+\frac{dM_7}{dR_7}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also normalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_7$) and the other derived from the non-delayed or "immediate" processes($I_7$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_7=\frac{N_7 G_7}{2M_7}\frac{d(7nF_7)}{dR_7}(\geq 0) \quad (7a)$$

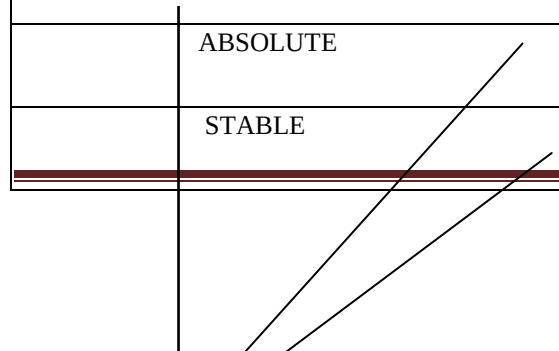
$$I_7=-\frac{N_7 G_7}{2M_7}\frac{d(7nM_7}{dR_7}(\geq 0) \quad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7>L_7-I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

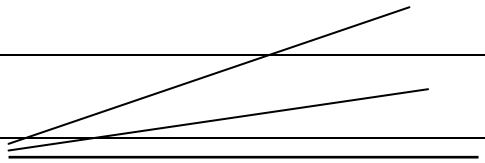
$$P_7>L_7-I_7, \quad \text{when} \quad P_7^2/4>I \quad (8a)$$

$$P_7>L_7-I_7+\frac{(I_7-P_7^2}{(L_7-I)} \quad \text{when} \quad P_7^2/4<I. \quad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,





**DELAYED PREDATOR CREDIT TYPES EFFECT WITH ITS WIDE RANGING RAMIFICATION**

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

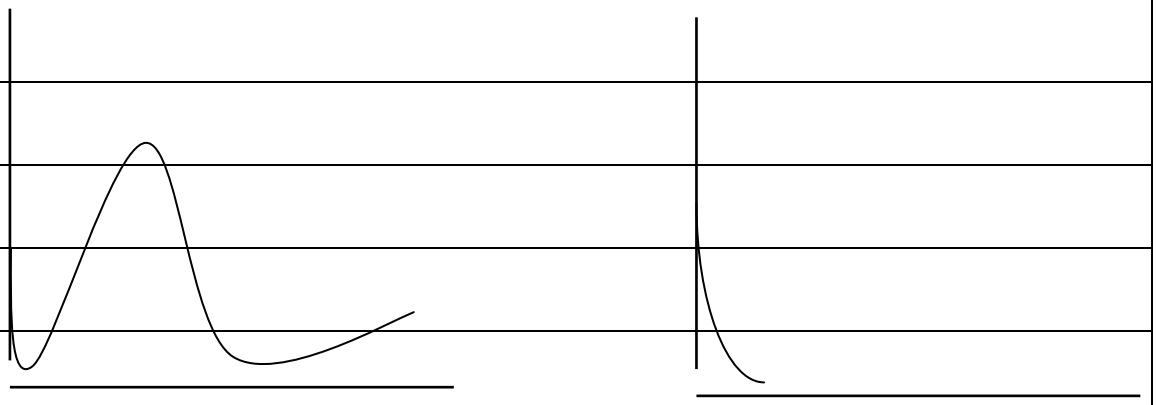
It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay

and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



**TIME DELAY**

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

To scores of problems this relationship is important. Of them some of the examples are as follows:

**New look at the galaxy-shaping power of black holes**

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| (769) | Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving. |
| (770) | "These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science. |
| (771) | "Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy." |
| (772) | This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return. |
| (773) | The remaining matter gets absorbed into the black hole, adding to (e) its mass. |
| (774) | The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne. |
| (775) | The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution. |
| (776) | Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form. |
| (777) | "Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara. |
| (778) | "It creates (eb) these bubbles, and the bubbles keep the plasma hot. |
| (779) | That's what prevents galaxies from (e) becoming even bigger than they are now." |
| (780) | Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space. |
| (781) | In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene. |
| (782) | Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve. |
| (783) | "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara. |
| (784) | "It's like nature's thermostat, which keeps these galaxies from (e) growing. |
| (785) | If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars." |

**A new look at the galaxy-shaping power of black holes**

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| (786) | An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the |
| (787) | McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said. |

(788) Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.

**Doomed satellite missions**

(789) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.

(790) At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.

(791) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.

(792) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.

(793) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.

(794) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people

(795) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."

(796) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.

(797) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(798) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.

(799) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.

(800) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.

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| (801) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses. |
| (802) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well. |
| (803) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma. |
| (804) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6. |
| (805) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus. |
| (806) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (807) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (808) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (809) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (810) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (811) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (812) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (813) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (814) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (815) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (816) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (817) We know dark matter exists only from (e) its gravitational effects. |
| (818) It holds together galaxies and (e&eb) galaxy clusters. |

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| (819) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (820) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |
| (821) They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (822) X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (823) They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (824) This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (825) An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |
| (826) Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (827) "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (828) The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (829) Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (830) Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (831) "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (832) Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (833) "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New |

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| York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (834) Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories. |
| (835) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen. |
| (836) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer. |
| (837) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants. |
| (838) These were treated as (=) "aware" and "unaware" trials, respectively. |
| (839) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it. |
| (840) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another. |
| (841) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion. |
| (842) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness. |
| (843) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said. |
| (844) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness." |
| (845) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner. |
| (846) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain. |
| (847) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said. |
| (848) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

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| (849) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter. |
| (850) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined. |
| (851) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions. |
| (852) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda. |
| (853) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general. |
| (854) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole. |
| (855) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes. |
| (856) Noda and his colleagues instead studied the change in emissions over (e&eb) time. |
| (857) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection. |
| (858) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1). |
| (859) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole. |
| (860) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved. |
| (861) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp |

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

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| (862) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking “palm tree” organization. |
| (863) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society |
| (864) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible? |
| (865) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem. |
| (866) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other. |
| (867) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible. |
| (868) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges. |
| (869) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century. |
| (870) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching. |
| (871) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom. |
| (872) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth. |

**Spheres on trees**

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| (873) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration. |
| (874) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement. |
| (875) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum. |
| (876) By visualizing the problem in this way, the researchers noticed that these particular graphs don’t (e) have lots of branches extending from lots of (e) other branches. |

(877) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure.

(878) The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network.

(879) As a result, it provides (eb) an important clue for finding the global minimum.

(880) It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum.

(881) This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large.

**Thompson's problem 2**

(882) By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society

(883) The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima.

**(884)** The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node.

**Implications for Thomson's and Smale's problems**

(885) Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before.

(886) "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around."

(887) As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem.

(888) "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view."

(889) In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges.

(890) From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear.

(891) "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK.

(892) "For larger numbers, there will likely be (=) multiple funnels.

(893) Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum."

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| (894) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem. |
| (895) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said. |
| (896) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org |

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t-T_7)F_7\big(R_7(t-T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t-T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) and the effective prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this

rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Look at the resemblance of String Theory and Abhinava's earlier theory and how the model could be applied: The world of manifestation and differentiated perceptions (Yikalpa) thus extends from one Centre to (eb) the next. **What is meant again is the realisation or the evolution of individual consciousness or general ledger to be (=)the cosmic general ledger is the sole criterion for the Satchidananda state where all dualities cease** Although it is never in fact divorced from (e) the subject who resides there, the ignorant fail (e) to grasp this fact and so, cut off from (e) the Centre, the world of objectivity becomes (=) for them the sphere of Maya. Bhagavatotpala quotes the Light of Consciousness amvitprakdsa): This ever pure experience (suddhanubhava) is variegated by (e) each form [revealed within it] Even so it remains unstained (nirmala) when moving to another. Just as a cloth which is naturally white, once dyed, cannot (e) change colour without [first] becoming white again, similarly the pure power of awareness, (citi) once coloured by (e) form, is pure [again] at the Centre where that form is (=) abandoned and from whence it proceeds to (e&eb) another. In his Essence of Vibration (Spandasarpdoha), K§emaraja explains that the rise and fall of every individual perception in the field of awareness is (=) a specific pulsation of consciousness. This is exactly the individual general ledger transaction lifted to the transaction in cosmic general ledger. During The Initial Instant Of Perception, T Consciousness Is Manifestly Apparent And The Yogi, Participating In Its Plenitude, Observes The Outer World Without Being Attached To Any Particular Divine Body And Sacred Circle Of The Senses Singling It Out From Any Other, Like A Man Who Observes A City From A High Mountain Peak.

He Sees The Outer World Reflected Within His Consciousness Free Of Thought-Constructs And So 'Stamps* The Outer On The Inner While Absorbing The Object And Means Of Knowledge In The Pure Subject Which Grasps Them As The Expansion Of His Own Nature. K§emaraja Says: By Penetrating Into Bhairavimudra, The Yogi Observes The Vast Totality Of Beings Rising From, And Dissolving Into, The Sky [Of Conscious- Ness], Like A Series Of Reflections Appearing And Disappearing Inside A Mirror. Through (e&eb) The Practice Of Bhairavimudra The Yogi Realises That He Is (=) The Substratum Consciousness (Adhisfhatr) Which Both Underlies (e&eb) And Is (=) The Essence Of All Things. He Discovers That Phenomena Have (e) No Independent Existence Apart From (e) Him And So Are, In This Sense, Void. At The Same Time, He Realises That Because (e) All Things Are Consciousness, They Are (=) Far From Unreal. He Views The Outer World Yet Sees (e) It Not. Beyond Both Voidness And Non-Voidness He Penetrates Into (e&eb) The Supreme Abode (Pararn Padam) Of Siva's Consciousness. [The Powers Of The Senses] Endowed With (e&eb) The Attributes Of The Great Union [Between Subject And (e&eb) Object] Whose Form Is The Wakening Of Man's Spiritual Potential (Kuntfalinf), Fill (e&eb) [With Consciousness] The Outer Clatter Of Diversity (Bhedatfambara) Born Of (e) Its Intense Power And Are Then Established In (eb) The Unobscured Abode Of The Void Of Consciousness To Shine [There] Eternally. Thus Residing Beyond (e) Being And Non-Being, The Sole Protector Of The Unity Which Is Tranquil And Expanding [Consciousness], Whose Glory Is All-Embracing And Form Unobscured, Is Called Bhairavimudra. Through (e&eb) The Practice Of Bhairavimudra, The Yogi Unites The Universal Vibration Of T Consciousness With (e&eb) The Individual Pulsation Of Objectivised 'This' Consciousness. The Two Aspects Of Consciousness Are Now In (eb) A State Of Equilibrium Like (=) The Two Pans Of An Evenly Weighed Balance And The Yogi Experiences (e&eb) The Pure Knowledge (Suddhavidya) That: 'I Alone Am (=)All Things'. Thus Becoming (=) The Master Of The Wheel Of Energies He Is (=) Free, Like Siva, To Create And Destroy. When [The Yogi] Is Well Established, Without (e) Wavering, Solely In The Integral Egoity Of His Authentic Nature, The Spanda Principle, And Is

Absorbed In Contemplation (Samavisfa), He Becomes (=) One With It The Doctrine Of Vibration (Tanmaya). Then . . . Dissolving And Creating The Universe By Means Of (er) His Introverted And Extroverted Absorption, He Destroys And Creates (e&eb) All Things Out Of Sarikara, His Innate Nature. [Thus] He Assumes The State Of The Universal Experiencer And Having Absorbed All That Is To Be Experienced From [The Grossest Level] — Earth — To (The Subtlest) — Siva — He Reaches (eb) The State Of The Supreme Subject By (e) Progressively Recognising [His Identity With Him]. Thus, Introverted And Extroverted Absorption Both Lead To (eb) The Recognition Of The Pulsation (Spanda) Of (e) One's Own Consciousness. This Is Where (e) Individual General Ledger Becomes (=) Cosmic General Ledger. At The Level Of Consciousness Corresponding To (e&eb) Siva's Basic State (Saryibhavdvastha), The Alternation From Inner To (e&eb) Outer Is Instantaneously Resolved Into (e&eb) The Vibration Of His Nature. When The Yogi Finally Comes To Be (=) Constantly Aware Of This Reality, His Enlightenment Is (=) Full And Perfect. Freed Of All Means (Anupaya) And Delighting In The Power Of His Bliss (Anandasakti), He Knows And Does Whatever He Pleases. The Yogi Seeking Self-Realisation Must Acquire (e) Mastery Over This Movement. Ksemaraja Stresses That The Doctrine Of Vibration Teaches That Liberation Can Only Be Achieved By (e) First Withdrawing All Sense Activity In Introverted Contemplation (Nimilanasamadhi) To (e) Then Experience The 'Great Expansion' (Mahdvikdsd) Of (e) Consciousness While (e) Recognising This To Be (=) A Spontaneous Process Within (eb) It. This Is Done Through (e&eb) The Practice Of Kramamudra. A Passage From The Now Lost Kramasiitra Explains: Although The Adept's Attention [May Be] Outwardly Directed, (eb) He Enjoys Contemplative Absorption Through (e&eb) The Introverted Aspect Of (e) Kramamudra. Initially He Turns Inward From (e)The Outside World And [Then] From Within [Himself] He Exits Into The Outer World Under The Influence Of His Absorption. There Is No Necessity Of Any Glorification Or Mortification Here. It Is Just Like Joining A Job And Retiring And Realising All That Happened Is Just By His Own Effort And Nothing Else. Thus The Sequence (Krama) In This Attitude (Mudra) [Ranges Through] Both Inner And Outer. The Yogi Must Pervade

The Surface Level Of Awareness (Yyutthana) With The Same Bliss He Experiences Plunging Into The Depths Of Con- Templative Absorption (Samadhi). Note Here That Samadhi Is A Death Like State And Still Being Matter And Information Stored You Experience Consciousness. Submerging Himself And Emerging Repeatedly From Samadhi, He Eventually Recognises That The Unity Of Consciousness Pervades Both States: The Best Of Yogi's, Who Has Achieved A State Of Complete Absorption Even When Risen From Meditation, [Inwardly] Vibrating Like A Drunkard In Blissful Inebriation From The After-Effects Of The Nectar Of Con- Templation, Sees All Things Dissolving In The Sky Of Consciousness Like A Cloud In The Autumn Sky. He Plunges Repeatedly within Himself Divine Body and Sacred Circle of the Senses and Becomes (=) Aware Of His Identity With (e&eb) Consciousness By (e) The Practice Of Introverted Contemplation. Thus Even When He Is Said To Have Risen From (e) Absorption, He Is (=) One With [His] Experience Of It. Ksemaraja Goes On To Explain That This Practice Is Called 'Mudra' Because (e) It Both Fills The Adept With (e&eb) Bliss (Mud) And Is Itself The Bliss Of (e) Consciousness. Moreover, It Dissolves (e) Away (Dra) All Bondage And 'Stamps' The Universe Of Experience With (e&eb)The Seal (Mudra) Of The Fourth State (Turiya) Of Enlightened Consciousness Beyond, And Including, The Three States Of Waking, Dreaming And Deep Sleep. It Is Called 'Krama' Because (e) It Is The Root Source Of (e) All Emanation And All Other Conscious Processes Which Succeed (e&eb) One Another In Ordered Sequence (Krama) And Is, (=)At The Same Time, Their Successive (Krama) Appearance As Well. By The Practice Of Kramamudra The Opposites Fuse And Siva And Sakti Unite. They Yogi Comes To Experience The Simultaneous Pervasion Of All The Lower, Grosser Categories Of Existence By The Higher And The Presence Of The Lower In The Higher. Commencing His Practice In A Low Form Of Bhairavimudra, The Yogi Conjoins The Outer With The Inner; Then, In Kramamudra, He Fills Both The Outer With (e&eb) The Inner And (e&eb) The Inner With (e&eb)The Outer. When He Achieves (eb) Perfection In This Two-Fold Movement, He Attains To (e) The Highest Form Of Bhairavimudra In Which The Two Merge Completely In (eb)The Experience Of (e) The Absolute (Anuttara), Free

Of (e) All Differentiation And (e&eb) Polarities. If He Fails To Maintain Awareness Of This State, He Again Falls Into (eb) Kramamudra Until He Has Finally Completely Merged All The Highest States In (e&eb) The Lower And The Lower In (e&eb) The Higher. He Then No Longer Needs To (e) Resort To (e) Any Means (Anupaya) To (e) Achieve Liberation. All He Says or Does Anything He Perceives or Thinks, Instantly Occasions In (eb) Him The Highest Level Of Consciousness. Thus The Fruit Of Bhairavimudra Is (*=)The Wonder (Camatkara) Or Amazement (Vismaya) That Overcomes (e)The Yogi When (e) He Reaches The Plane Of (e) Union (Yogabhumika) Y Xsl Where (e) All Opposites Merge In (e&eb)The Radiance Of The Great Light Of Consciousness. The Stanzas On Vibration Teach: How Can One Who, As If Astonished, Beholds (e) His Own Nature As (=) That Which Sustains [The Existence Of Everything] Be Subject To (e&eb) This Painful Round Of Transmigration (Kusrti)? The Yogi, Recognising His True Nature To Be (=)The Supreme Subject, Is Astonished To (e) Suddenly Discover That The Individual He Thought He Was, Caught Up (e&eb) In The Trammels Of Thought And Living In A World Enmeshed In (e&eb) The Web Of Time And Space, Does Not Really Exist At All. He Experiences A 'Turning About' (Paravrtti) In (eb)The Deepest Seat Of The Doctrine Of Vibration Consciousness As He Penetrates His True Nature. The Sudden Eruption Of This Intuition (Pratibha) Arouses In Him A Cry Of Amazement As He Transcends All Thought-Constructs And, Perfectly Absorbed In His Own Nature, Is Liberated. The Path To Liberation Essentially, Spanda Doctrine Is Concerned With Two Matters. The First Is To Impart To Those Who Are Fit To Receive The Teachings A Deeper Understanding Of The Ultimate Goal Of Life (Upeya). When We Have Understood What Truly Benefits Us And Is Worth Attaining And What, On The Contrary, Is Of No Real Value But Stands In The Way Of This Attainment, We Can Begin To Make Progress Towards Our Goal. This Is Spanda Doctrine's Second Concern, Namely, To Show The Way In Which We Can Develop Spiritually Through Siva's Grace And The Right Application Of The Means To Realisation That It Teaches. When Both These Aspects Of The Teaching Have Been Correctly Understood And Applied, The Spanda Yogi Achieves A Clear And Permanent

Realisation Of His Goal And Is Liberated, Thus Fulfilling The Ultimate Aim Of The Teaching. The Doctrine Of Vibration Is Not Meant For The Spiritually Dull. It Is Not For The Worldly Whose Consciousness, Clouded By Ignorance, Is As If Dreaming, Even During The Waking State Of Daily Life, The Dream Of Its Own Thought-Constructs. The Teachings Are Meant For Those Who Are Awake (Prabuddha), Those Who, Full Of Faith And Reverence, Are Always Alert And Intent On Discerning The True Nature Of Ultimate Reality. This Reality Is Understood In Three Basic Ways. The First Is Purely Transcendental. The Stanzas Choose This Aspect As The One Which Formally Defines It Most Specifically. Ultimate Reality Transcends All The Opposites, Including Subject And Object. This Does Not Mean, However, That It Is An Unconscious Void, A Mere Absence Of All Existence. In Fact, This Negative Characterization Of Reality (Which Includes Also A Denial Of All That Is Unconscious) Implies A Positive Immanence In Which The Opposites Are United In The Oneness Of Pure Consciousness That Is Equally Siva And Spanda, His Universal Activity. These Two Seemingly Contrasting Aspects Are Reconciled In The Third, Namely, Reality Understood As The Essential Nature Of All Things. Although Universal And Everywhere The Same, It Is Understood To Be The Essential And Specific Nature Of Each Existent As Its 'Own Nature' (Svabhava).

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such

a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

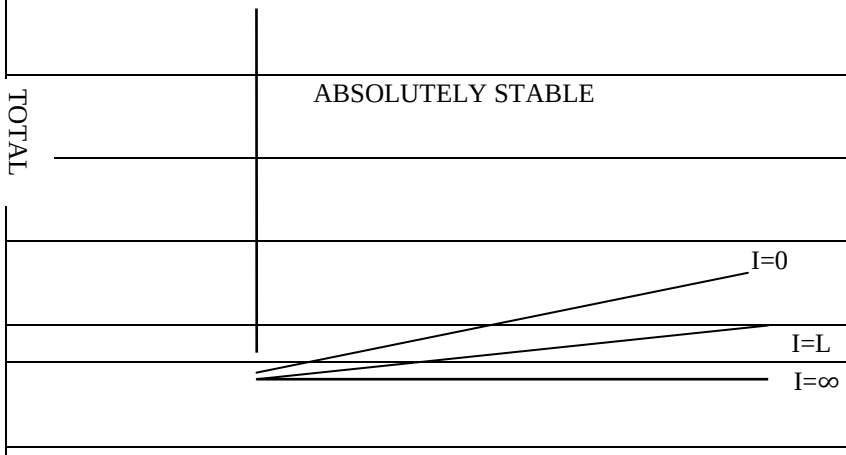
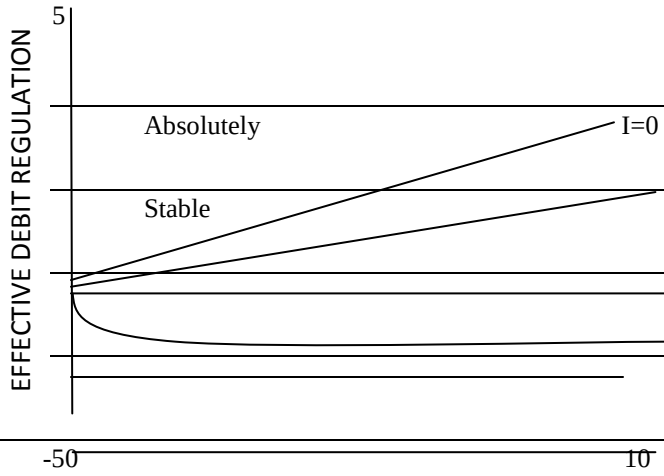



FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

This equation is of paramount importance in the following cases: According to the principles of geometric quantum mechanics, the physical characteristics of a given quantum system can be represented by geometrical features that are preferentially identified in this complex manifold. Here we construct a number of examples of such features as they arise in the state spaces for spin, spin 1, spin and spin 2 systems, and for pairs of spin systems. A study is then undertaken on the geometry of entangled states. A locally invariant measure is assigned to the degree of entanglement of a given state for a general multi-particle system, and the properties of this measure are analysed for the entangled states of a pair of spin particles. With the specification of a quantum Hamiltonian, the resulting Schrödinger trajectories induce an isometry of the Fubini–Study manifold, and hence also an isometry of each of the energy surfaces generated by level values of the expectation of the Hamiltonian. For a generic quantum evolution,

the corresponding Killing trajectory is quasiergodic on a toroidal subspace of the energy surface through the initial state. When a dynamical trajectory is lifted orthogonally to Hilbert space, it induces a geometric phase shift on the wave function. The uncertainty of an observable in a given state is the length of the gradient vector of the level surface of the expectation of the observable in that state; a fact that allows us to calculate higher orders corrections to the Heisenberg relations. A general mixed state is determined by a probability density function on the state space, for which the associated first moment is the density matrix. The advantage of a general state is in its applicability in various attempts to go beyond the standard quantum theory, some of which admit a natural phase-space characterization. **Journal of Geometry and Physics Volume 38, Issue 1, April 2001, Pages 19–53 Geometric quantum mechanics Dorje C. Brody Lane P. Hughston DOI: 10.1016/S0393-0440(00)00052-8** Expansion (unmesa) of [consciousness] or the creative intuition {pratibha) is (=) [experienced] in the interval which divides two [moments] of differentiated perception (vikalpa) It is here that they arise and disappear. The Sastras and Agamas proclaim with reasoned argument that it is (=) free of thought-constructs {Nirvikalpa) and precedes (e) Path to Liberation all mental representations of any object. None can deny that a gap exists between perceptions insofar as two moments of thought are (=) invariably divided. This [gap] is (=) the undifferentiated unity of all the countless manifestations. Similarly, in the outer more objective sphere, where change consists (e) of the alterations in the configurations of manifest appearances (Abhasa\ the transition from one to (e&eb) another corresponds to (e&eb) a phase of pure luminosity that marks (eb) the beginning of one form and the end of another

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^{\infty} \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7 \left( U_7 - V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7 \right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$then given the constraints of $f(T_7), (A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7 \lambda_7 e^{-T_7 \lambda_7} + [(U_7 - V_7) M_7 + Y_7] - [((U_7 -$$

V 7M 7−]e−T 7λ 7=0

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$\boldsymbol{Z}_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7) \cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7) \sin(Z_7 T_7) = 0$$

The warrant of Truth is ever elusive when we deconstruct the foundational justifications of our convenient notions about the way the world works. It is easy to confuse what is actually the creation of beliefs with the "discovery of Truth," a common goal of science and theology. Science offers no ontological Scientific Picture of the World but substitutes a number of applied theories as meta-theory. It's close, powerful and useful: it 'works' and allows use to do work. But the still mysterious "true" metatheory may reveal the basis of consciousness. But who knows when we will lift that veil? Permeating the living reality of our culture are certain contagious notions, fads and trends that have the ability to influence the way we think about the nature of Reality and ourselves. Some of them are toxic and can consume you; others are just whims of pop culture infectious and benign as the common cold. **Jung** described the concept of psychic contagion by certain archetypal forces inherent in the human psyche, which manifest in our spiritual lives and belief systems. He spoke of both conscious and unconscious contamination. Notions like this range from simple superstitions to scientific concepts, to urban myths. Notions sweep through our culture and insinuate themselves within its fabric, as fads, whether they are "real" or not, they can be influential. They mood alter us, make us feel we belong. An analysis of these notions is useful in distinguishing a common human phenomenon from any potential "alien influence" which may or may not be exerted on us from an unknown source. **Strange Attractors: Transference, Holography, and an Archetype Burke, J. (2003).**

**Strange Attractors: Transference, Holography, and an Archetype (Doctoral dissertation, Pacifica Graduate Institute, 2003).** In consideration to the Nobel Proze for Chemistry this year(2014), a stability nalysis of the proble in to single molexues would be a contemporaneous question that has to be taken cognisace of and given importance to . Determination of molecular structure by geometry optimization became routine only after efficient methods for calculating the first derivatives of the energy with respect to all atomic coordinates became available. Evaluation of the related second derivatives allows the prediction of vibrational frequencies if harmonic motion is estimated. More importantly, it allows for the characterization of stationary points. The frequencies are related to the eigenvalues of the Hessian matrix, which contains second derivatives. If the eigenvalues are all positive, then the frequencies are all real and the stationary point is a local minimum. If one eigenvalue is negative (i.e., an imaginary frequency), then the stationary point is a transition structure. If more than one eigenvalue is negative, then the stationary point is a more complex one, and is usually of little interest. When one of these is found, it is necessary to move the search away from it if the experimenter is looking solely for local minima and transition structures. The total energy is determined by approximate solutions of the time-dependent Schrödinger equation, usually with no relativistic terms included, and by making use of the Born–Oppenheimer approximation, which allows for the separation of electronic and nuclear motions, thereby simplifying the Schrödinger equation. This leads to the evaluation of the total energy as a sum of the electronic energy at fixed nuclei positions and the repulsion energy of the nuclei. Notable exceptions are certain approaches called direct quantum chemistry, which treat electrons and nuclei on a common footing. Density functional methods and semi-empirical methods are variants on the major theme. **Orlando Alvarez** discusses Polyakov's quantization of the string in the presence of a boundary allowing for an arbitrary topology for the world sheet. In addition to the dynamical conformal factor discovered by Polyakov, there are a finite

number of new degrees of freedom if the surface is more complicated than a sphere or a disc. The quantization of the Liouville theory in an arbitrary topology is also discussed. A one-loop calculation shows that the model is renormalizable if one performs a mass renormalization and an additive field renormalization. The renormalization group equations have a perturbative infrared unstable fixed point in all topologies. Copyright © 1983 Published by Elsevier B.V. **Nuclear Physics B Volume 216, Issue 1, 25 April 1983, Pages 125–184 Theory of strings with boundaries: Fluctuations, topology and quantum geometry Orlando Alvarez DOI: 10.1016/0550-3213(83)90490-X** Quantum geometry (the modern loop quantum gravity involving graphs and spin-networks instead of the loops) provides microscopic degrees of freedom that account for black-hole entropy. However, the procedure for state counting used in the literature contains an error and the number of the relevant horizon states is underestimated. In our paper a correct method of counting is presented by **Marcin Domagala and Jerzy Lewandowski.** Results lead to a revision of the literature of the subject. It turns out that the contribution of spins greater than 1/2 to the entropy is not negligible. Hence, the value of the Barbero–Immirzi parameter involved in the spectra of all the geometric and physical operators in this theory is different than previously derived. Also, the conjectured relation between quantum geometry and the black-hole quasi-normal modes should be understood again. **Marcin Domagala and Jerzy Lewandowski 2004 Class. Quantum Grav 21 5233 doi:10.1088/0264-9381/21/22/014 Black-hole entropy from quantum geometry** Quantum geometry predicts that a universe evolves through an inflationary phase at small volume before exiting gracefully into a standard Friedmann phase. This does not require the introduction of additional matter fields with ad hoc potentials; rather, it occurs because of a quantum gravity modification of the kinetic part of ordinary matter Hamiltonians. Authors draw the cognizance of application of the same mechanism that can explain why the present day cosmological acceleration is so tiny. DOI: http://dx.doi.org/10.1103/PhysRevLett.89.261301

**Inflation from Quantum Geometry Phys. Rev. Lett 89, 261301 – Published 12 December 2002 Martin Bojowald** The loop quantum cosmology of the closed isotropic model is studied with special emphasis on a comparison with traditional results obtained in the Wheeler-DeWitt approach by **Martin Bojowald** .Investigation includes the relation of the dynamical initial conditions to boundary conditions such as the no-boundary or the tunneling proposal and a discussion of inflation from quantum cosmology. **DOI: http://dx.doi.org/10.1103/PhysRevD.67.124023 Loop quantum cosmology, boundary proposals, and inflation Phys. Rev. D 67, 124023 – Published 23 June 2003 Martin Bojowald and Kevin Vandersloot** In loop quantum cosmology, the universe avoids a big bang singularity and undergoes an early and short super-inflation phase. During super-inflation, non-perturbative quantum corrections to the dynamics drive an inflaton field up its potential hill, thus setting the initial conditions for standard inflation. **Shinji Tsujikawa et al** show that this effect can raise the inflaton high enough to achieve sufficient e-foldings in the standard inflation era. They also analyse the cosmological perturbations generated when slow-roll is **violated after super-inflation** and show that loop quantum effects can in principle leave an **indirect signature** on the largest scales in the CMB, with some loss of power and running of the spectral index. **Shinji Tsujikawa et al 2004 Class. Quantum Grav 21 5767 doi:10.1088/0264-9381/21/24/006 Loop quantum gravity effects on inflation and the CMB Quantum geometry from quantum information** is the study of **Florian Girelli and Etera R Livine .**Loop quantum gravity defines the quantum states of space geometry as spin networks and describes their evolution in time. We reformulate spin networks in terms of harmonic oscillators and show how the holographic degrees of freedom of the theory are described as matrix models. This allows us to make a link with non-commutative geometry and to look at the issue of the semi-classical limit of loop quantum gravity from a new perspective. This work is thought of as part of a bigger project of describing quantum geometry in quantum information terms.

**Florian Girelli and Etera R Livine 2005 Class. Quantum Grav 22 3295 doi:10.1088/0264-9381/22/16/011 Reconstructing quantum geometry from quantum information: spin networks as harmonic oscillators.** Quantum geometric mechanics is the topic of study by **Dorje C. Brody Lane P. Hughston** establishing a clear relation between the quantum mechanics and quantum geometry and also the entangled states. The manifold of pure quantum states can be regarded as a complex projective space endowed with the unitary-invariant Fubini–Study metric. Note how differential gheometry concepts with Nunney's model could be entangle so that it is feasibly and viably applied to various type of model including consciousness (See Doctrine of Vibration)

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $\boldsymbol{Z}_7^2$ and $Z_7$in (A2.2):

$W_7^2 (X_7^2$-3 $M_7\ A_7\ X_7$- $M_7^2 (X_7$- $Y_7) + 2\ M_7^3\ A_7$+2 $A_7^2\ M_7^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_7$=1-cos ($Z_7 T_7$).

In the absence of a delay the system is stable for $A_7 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7 > (X_7 - Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$A_7 > [\ X_7 - Y_7/2\ M_7$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_7, A_7) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_7 \le 2$ *is* of interest First consider the extreme values

$$Q_7(0, A_7) = A_7^2 (X_7 - Y_7)$$

$$Q_7(2, A_7) = (A_7 + 2M_7)^2 (2A_7 M_7 - (X_7 - Y_7))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7 M_7^2 - 3M_7 X_7\right) + W_7\left(3M_7 A_7^2 - 2A_7 X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7 + Y_7\right)$$

$$=2M_7^2 W_7^2 (M_7 + A_7) + M_7 A_7^2 W_7 + (2A_7 M_7 - X_7 + Y_7)(A_7 W_7 + M_7 W_7^2$$

$$+ (2 - W_7)((2M_7 X_7 + M_7 Y_7) W_7 + A_7 (X_7 + Y_7)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $W_7$touches the $W$–axis only at its minimum value in which case (A2.9) shows that

further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$has no minimum in the range 0≤ $W_7$ ≤2.Defining

$Q_7(W_7, A_7) = a_7 W_7^2 + c_7$ *where* $a_7\, b_7$ *and* $c_7$ Are functions of $A_7$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7 \geq b_7$$

And

$$b_7^2 \geq 4a_7 c_7$$

$$Case(i): Y_7 = 0$$

*Let* $Y_7 = 0$ And try the substitution A= $X_7 = /2M_7$ *in* (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_7 X_7 /$ $2M_7$) is positive for *W<2. Since* $Q_7$ increases with $A_7$(see A2.9), increasing $A_7$will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

$Case\ (ii): Y_7 = y_7\ X_7$ ,

$Let\ Y_7 = y_7\ X_7$And try the substitution $A_7 = X_7/(2 + A_7\ y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + n_7 y_7)$ we have

$$W_7^2\ y_7\left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2\left(2 - n_7 + n_7 y_7\right)\right]$$

$$-W_7\ X_7\left[\frac{X_7(1 + n_7 y_7)}{M_7^2(2 + n_7 y_7)^2} + 2\ y_7\left(n_7 - 1\right) + \frac{X_7^2(7 + y_7)}{M_7^2(2 + n_7 y_7)} = 0.\ (A2.12)\right.$$

Since $n_7$ and $A_7$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7), the\ condition (A2.77?) can\ be\ written\ as$

$[\frac{X_7(E_7-7)}{E_7^2}]^2[\frac{X_7^2}{M_7^2} - 4\ y_7\ E_7^2\frac{X_7}{M_7^2} + 4y_7(n_7 - 2)E_7^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$2\ E_7[y_7 E_7 + (y_7\ E_7(7 + y_7)(E_7 - n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - y_7 E_7(7 + y_7)(E_7 - n_7))^{7/2}]$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0 $W_7$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7-7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2-2y_7\frac{(E_7+n_7^2y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right)-4y_7(E_7-n_7)$$

And has the solution

$$\frac{X_7}{M_7^2}>\frac{E_7}{(E_7-7)}[y_7(E_7+n_7^2y_7)+[y_7^2(E_7+n_7^2y_7)^2+4y_7(E_7-7)(E_7-n_7)]^{7/2}]$$

Given $E_7>n_7$this lower limit to $X_7/M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7/M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$in the appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7<n_7$ $(ie.,2A_7M_7^2<X_7-Y_7$)then (A2.7)and (A2.8)show that $Q_7$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2}>2E_7[y_7E_7-(y_7E_7(7+y_7)(E_7-n_7)]^{7/2}]$$

And

$$E_7>n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2 > 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.

If $A_7^2 < 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.+ $(Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$ and $X_7$ and $Y_7$ by $W_7^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7\lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]\, e^{2T_7\lambda_7} + M_7 V_7\, e^{2T_7\lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i}\, Z_7$ and separate

(A3.1)into real and imaginary parts

$$Z_7^2 + (M_7 + V_7) Z_7 \sin(Z_7 T_7) - (U_7 M_7 + Y_7) + (M_7(U_7 + V_7) - X_7)$$

$$\mathrm{X} \cos(Z_7 T_7) - 2M_7 V_7 \cos^2(Z_7 T_7) + M_7 V_7 = 0$$

$$(M_7 + U_7) Z_7 - (M_7 + V_7) Z_7 \cos(Z_7 T_7) + (M_7(U_7 + V_7) - X)$$

$$\mathrm{X} \sin((Z_7 T_7) - M_7 V_7 \sin(Z_7 T_7) = 0$$

Using (A.3.3) to eliminate $Z_7$and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - cos\, Z_7\ T_7$and $A_7 = U_7 - V_7$we obtain

$$4\,M_7^2 V_7^2 W_7^4 + 2\,M_7 V_7\ W_7^3[(M_7 - V_7)^2 + A_7(M_7 + V_7) + 2(A_7 M_7 - X_7)]$$

$$+ W_7^2[M_7 - V_7)^2(2\,A_7\ M_7 + Y_7 - X_7) + (2A_7\ M_7 - X_7)^2 - A_7(M_7 - V_7)(2A_7\ M_7 - X_7)$$

$$+ M_7 V_7\,(A_7^2 + 4(X_7 + Y_7))]$$

$$+ W_7[A_7^2 + 2X_7)(2\,A_7\ M_7 - X_7) + 2A_7 Y_7(M_7 + V_7) - M_7 A_7^3] + A_7^2(X_7\ M_7 - Y_7) = 0$$

Being a quartic in $W_7$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7 > 0$$

$$and$$

$$C_7 > (X_7 - Y_7)/2\,M_7$$

$recalling\ that\ A_7 = U_7$and defining $C_7 = U_7 + V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0, U_7 V_7 = A_7^2(X_7 + Y_7)$$

$$Q_7\left(2, U_7 V_7\right) = \left(C_7 + 2M_7^2\right)^2 (2C_7 M_7 + Y_7 - X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

## Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

**Module numbered five**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This is the condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference

between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \quad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t - T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\left(R(t)\right)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

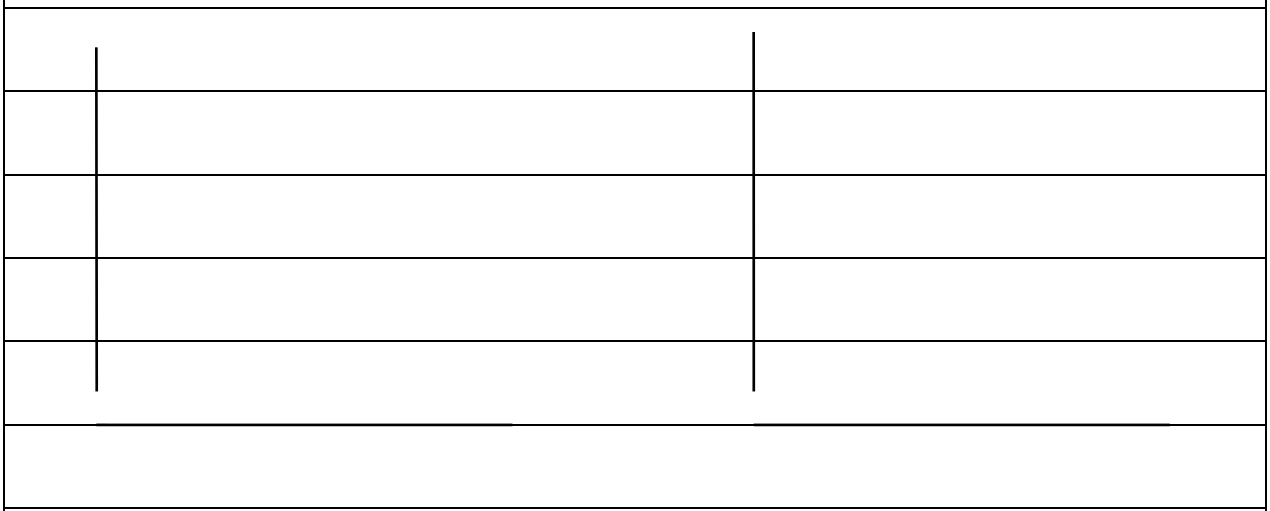
Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time

delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the



TIME DELAY

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period,

which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

As it is usually expressed outside of phenomenological texts, to undergo a conscious experience necessarily means that there is something it is like for the subject to have that experience (Nagel 1974; Searle 1992). This is obviously true of bodily sensations like pain. But it is also the case for perceptual experiences, experiences of desiring, feeling, and thinking. There is something it is like to taste chocolate, and this is different from what it is like to remember what it is like to taste chocolate, or to smell vanilla, to run, to stand still, to feel envious, nervous, depressed or happy, or to entertain an abstract belief. Yet, at the same time, as I live through these differences, there is something experiential that is, in some sense, the same, namely, their distinct first-personal character. All the experiences are characterized by a quality of mineness or for-me-ness, the fact that it is I who am having these experiences. All the experiences are given (at least tacitly) as my experiences, as experiences I am undergoing or living through. All of

this suggests that first-person experience presents me with an immediate and non-observational access to myself, and that consequently (phenomenal) consciousness consequently entails a (minimal) form of self-consciousness. To put it differently, unless a mental process is pre-reflectively self-conscious there will be nothing it is like to undergo the process, and it therefore cannot be a phenomenally conscious process. The mineness in question is not a quality like being scarlet, sour or soft. It doesn't refer to a specific experiential content, to a specific what; nor does it refer to the diachronic or synchronic sum of such content, or to some other relation that might obtain between the contents in question. Rather, it refers to the distinct givenness or the how it feels of experience. It refers to the first-personal presence or character of experience. It refers to the fact that the experiences I am living through are given differently (but not necessarily better) to me than to anybody else. It could consequently be claimed that anybody who denies the for-me-ness of experience simply fails to recognize an essential constitutive aspect of experience. Such a denial would be tantamount to a denial of the first-person perspective. It would entail the view that my own mind is either not given to me at all — I would be mind- or self-blind — or is presented to me in exactly the same way as the minds of others. There are also lines of argumentation in contemporary analytical philosophy of mind that is close to and consistent with the phenomenological conception of pre-reflective self-awareness. Alvin Goldman provides an example: [Consider] the case of thinking about x or attending to x. In the process of thinking about x there is already an implicit awareness that one is thinking about x. There is no need for reflection here, for taking a step back from thinking about x in order to examine it…When we are thinking about x, the mind is focused on x, not on our thinking of x. Nevertheless, the process of thinking about x carries with it a non-reflective self-awareness (Goldman 1970, 96). A similar view has been defended by Owen Flanagan, who not only argues that consciousness involves self-consciousness in the weak sense that there is something it is like for the subject to have the experience, but also speaks of the low-level self-consciousness involved in experiencing my experiences as mine (Flanagan 1992, 194). As Flanagan quite correctly points out, this primary type of self-consciousness should not be confused with the

much stronger notion of self-consciousness that is in play when we are thinking about our own narrative self. The latter form of reflective self-consciousness presupposes both conceptual knowledge and narrative competence. It requires maturation and socialization, and the ability to access and issue reports about the states, traits, dispositions that make one the person one is. To claim that every kind of self-consciousness is conceptual is overly cognitive. Bermúdez (1998), to mention one further philosopher in the analytic tradition, argues that there are a variety of **nonconceptual forms of self-consciousness** that are "logically and ontogenetically more primitive than the higher forms of self-consciousness that are usually the focus of philosophical debate" (1998, 274; also see Poellner 2003). This growing consensus across philosophical studies supports the phenomenological view of pre-reflective self-consciousness. (References: Stanford Encyclopedia, Post Modersnist Theories, Internet Encyclopedia of Philosophy) Korzybski pointed out that there are two ways to slice easily through life; to believe everything or to doubt everything. Both ways save us from thinking. So, from time to time you have to give yourself a Virtual Reality Check. The hallmark of postmodern philosophy has been disbelief or skepticism of all "metanarratives," or translations of reality. Postmodernism has even turned its profound skepticism on such important humanist concepts as "objective truth" and reason.Yet, for a deconstructionist postmodern society, individually we are still riddled with superstition and gullibility, and open to manipulation through our belief systems as any politician, philosopher, clergy, or salesperson will attest. Further, most people are painfully naive when it comes to even the simplest scientific understanding. Most of us don't have a clue about the fundamental nature of physical reality or our own psychological nature, and our ability to be fooled by our senses and mind, or methods of social and technological persuasion. We may think we know, but the penetration is deeper than we dare imagine. Nature rejects the naiveté that seeks absolute truth. We are beginning to realize, individually and culturally, that "realities" are all human constructions. The task becomes one of "catching ourselves in the act" of creating our own "reality" from the flow of events. Human truth is always an engagement of mind with experience.. We don't need to fear the collapse of our

personalistic belief system (the "box" we live in), nor our belief in absolute truth. **A strong desire to engage in the "quest for uncertainty" complements our anxiety that perhaps there is no absolute, objective ground to reality**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t-T_8)F_8(R_8(t-T_8) - N_8(t)M_8((R_8(t) \qquad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8} + \frac{N_8 dG_8}{dR_8} - \frac{dB_8}{dR_8} > \frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8} + \frac{dM_8}{dR_8}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_8 = \frac{N_8 G_8}{2M_8}\frac{d(8nF_8)}{dR_8} (\geq 0) \quad (8a)$$

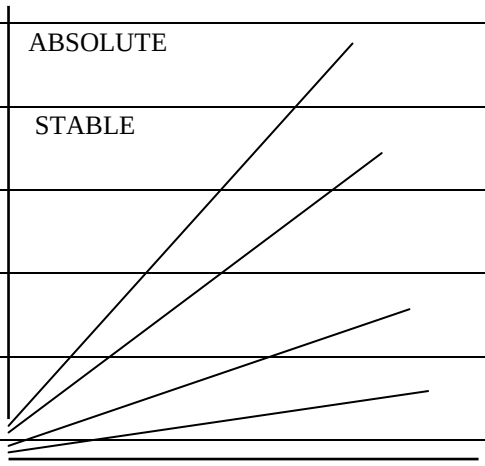
$$I_8 = -\frac{N_8 G_8}{2M_8}\frac{d(8nM_8}{dR_8} (\geq 0) \quad (8b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8 > L_8 - I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_8 > L_8 - I_8$, when $P_8^2/4 > I$ (8a)

$P_8 > L_8 - I_8 + \frac{(I_8 - P_8^2}{(L_8 - I)}$ when $P_8^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

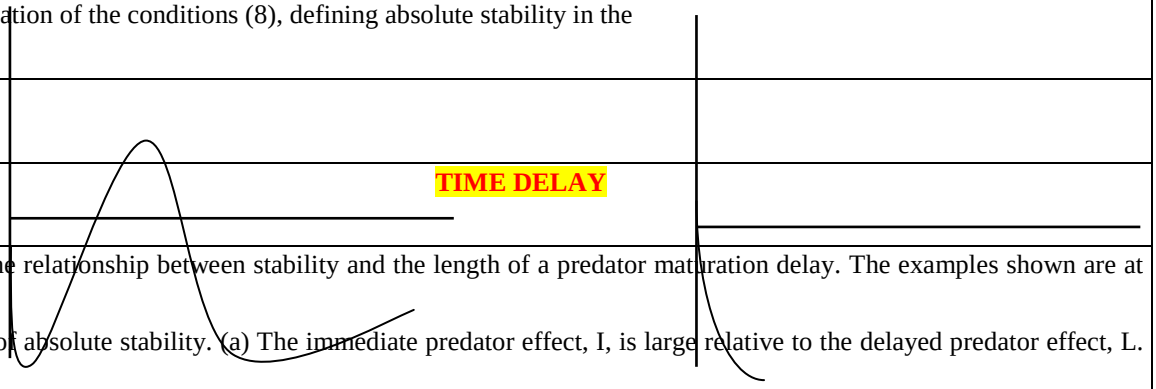



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this

rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior , we must note that the model could be applied to psychological systems too, albeit we are still not sure of the fact how the wave function collapses thereby due to perception.

Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility. He privatizes the sublime, having learned from the fate of his predecessors that the public can never be more than beautiful."(11) Notice is made here of what is indeed the telling feature. Levinas is praised, for example, for his alarm when "the social will [is] sought in an ideal of fusion . . . the subject losing himself in a collective representation, in a common ideal. . . . It is the collectivity which says 'us,' and which, turned toward the intelligible sun, toward the truth, experience, the other at his side and not face to face with him. The tellingly titled D'un ton apocalyptique contains passages that are particularly excited: These people situate themselves outside the ordinary, but they have in common this: they describe themselves as having an immediate and intuitive relation with mystery. And they want to attract, to seduce, and lead others to the mystery, through mystery. This agogic function of the leader of men, of the duce, of the Führer, of the leader, places him above the crowd that he manipulates with the aid of a small numbers of adepts joined together in a sect with a secret language, a clique or a small party with its

ritualized practices. The mystifiers pretend to have exclusive access to the privilege of a secret mystery. . . . The revelation or the unveiling of the secret is something that they jealously reserve for themselves. Jealousy is here a major characteristic. There is not narcissism and non-narcissism; there are narcissisms that are more or less comprehensive, generous, open, extended. What is called non-narcissism is in general but the economy of a much more welcoming, hospitable narcissism, one that is much more open to the experience of the other as other. I believe that without a movement of narcissistic reappropriation, the relation to the other would be absolutely destroyed; it would be destroyed in advance. The relation to the other--even if it remains asymmetrical, open, without reapporpriation--must trace a movement of reappropriation in the image of oneself for love to be possible, for example. Love is narcissistic. Beyond that, there are little narcissisms, there are big narcissisms. . . .(23) **Justice of the Pieces Deconstruction as a Social Psychology Douglas Collins** Sabine Wilke cogently argues for the rise of the aesthetic dimension in philosophical writing by showing how the metaphorical and poetic style of Heidegger, the principles of phrase configurations with Adorno, or the topographic writing of derrida replace through their very form something what traditionally would have been subject to discursive exposure The concern shifts from telling to showing, showing first and foremost that something remains untold. An exemplary book in this aesthetic drive is Avital Ronell's Telephone Book where the playful layout and explicit break with traditional design interrupt the sequential logic of the paragraphs thereby performing the very argument at stake: open up new circuits of meaning and signification. The engagement with form and publishing formats is only consequential since media theory as exposed by Walter J. Ong, Eric A. Havelock, Parry Miman or Marshall McLuhan besides many others made it clear that different communication technologies condition the mode of thought of a given culture The fractals of Benoit Mandelbrot or the butterfly shaped Lorenz attractor are the popular icons of this turn in the sciences. For communication studies and philosophy, Vilém Flusser was one of the pioneer thinkers reflecting the rise of what he called “techno-images.” The photograph is the paradigmatic case which condenses the whole crises of Western

thought and culture on its grainy surface. It subverts the linear logic of writing.6 Flusser's central diagnosis detects a shift from the grand narratives associated with the literate mind and Modernity towards creating and envisioning the world through image based practices. Writing himself in a succinct, visually inspired style, Flusser contents that the web of meaning, spun throughout the centuries by our textual practices, has become transparent. We discover the void between the letters through which the world escapes: "We are alienated from the world of text, precisely because we see through them as our own product."7 As a consequence we are confronted with a world of dissociated facts and information bits. Flusser hails this fragmented world as the advent of the zero dimensional point-universe that will provide the new raw material for our image based practices. Techno-images are nothing else than a conscious, inherently artistic drawing-together of point-elements as exemplified first and foremost by the silver-grains of the photograph but also by more recent pixel based techno-images. Similar to blueprints, designs or statistical curves, the weather map for example visualizes and contracts a tremendous amount of meteorological data from widely distributed stations into a single image of high- and low pressure

areas, temperatures, wind speeds, etc. Techno-images "compute" in this sense the universe of unimaginable, abstract information-bits into a concrete surface – the image. Echoing Paul Feyerabend, these image practices have to be understood as creative practices where the truth of scientific models has been replaced by its aesthetic quality **Aesthetics, Writing, Networked Computers Jörg Muller Dissertation Submitted to the Division of Media and Communications of The European Graduate School in Candidacy for the Degree of Doctor of Philosophy**

**Lothar Schäfer (1) and Sisir Roy (2) in** Quantum Reality, the Importance of Consciousness in the Universe, the Discovery of a Non-Empirical Realm of Physical Reality, and the Convergence with Ancient Traditions of Indian and Western Philosophy investigate metaphysical aspects that include 1) the discovery of a non-empirical part of physical reality in a realm of potentiality; 2) the emanation of the empirical world out of a realm of non-material forms; 3) the discovery that the nature of physical reality is that of an indivisible Wholeness – the One; and 4)

panpsychism: the possibility that the One is aware of its processes like a Cosmic Consciousness. The convergence of powerful traditions of seemingly disparate cultures is particularly important to point out in the present process of globalization, when a unifying view is needed to avoid controversy and conflict. To quote them in extenso : This is the nonlocality of the quantum world. Menas Kafatos and Robert Nadeau (1990) have drawn a remarkable conclusion from this phenomenon: If reality is nonlocal, the nature of the universe is that of an undivided wholeness. Because our consciousness has emerged from this wholeness and is part of it, it is possible to conclude that an element of consciousness is active in the universe: a Cosmic Consciousness. Pre-reflective self-consciousness is pre-reflective in the sense that (1) it is an awareness we have before we do any reflecting on our experience; (2) it is an implicit and first-order awareness rather than an explicit or higher-order form of self-consciousness. Indeed, an explicit reflective self-consciousness is possible only because there is a pre-reflective self-awareness that is an on-going and more primary self-consciousness. Although phenomenologists do not always agree on important questions about method, focus, or even whether there is an ego or self, they are in close to unanimous agreement about the idea that the experiential dimension always involves such an implicit pre-reflective self-awareness. In line with Edmund Husserl (1959, 189, 412), who maintains that consciousness always involves a self-appearance (Für-sich-selbst-erscheinens), and in agreement with Michel Henry (1963, 1965), who notes that experience is always self-manifesting, and with Maurice Merleau-Ponty who states that consciousness is always given to itself and that the word ‘consciousness’ has no meaning independently of this self-givenness (Merleau-Ponty 1945, 488), Jean-Paul Sartre writes that pre-reflective self-consciousness is not simply a quality added to the experience, an accessory; rather, it constitutes the very mode of being of the experience: This self-consciousness we ought to consider not as a new consciousness, but as the only mode of existence which is possible for a consciousness of something (Sartre 1943, 20 [1956, liv]). The notion of pre-reflective self-awareness is related to the idea that

experiences have a subjective 'feel' to them, a certain (phenomenal) quality of 'what it is like' or what it 'feels' like to have them.

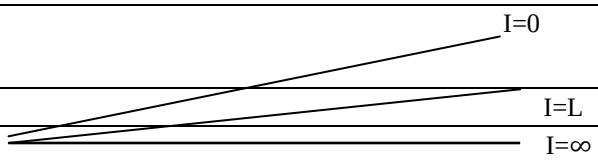
Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

TOTAL

ABSOLUTELY STABLE

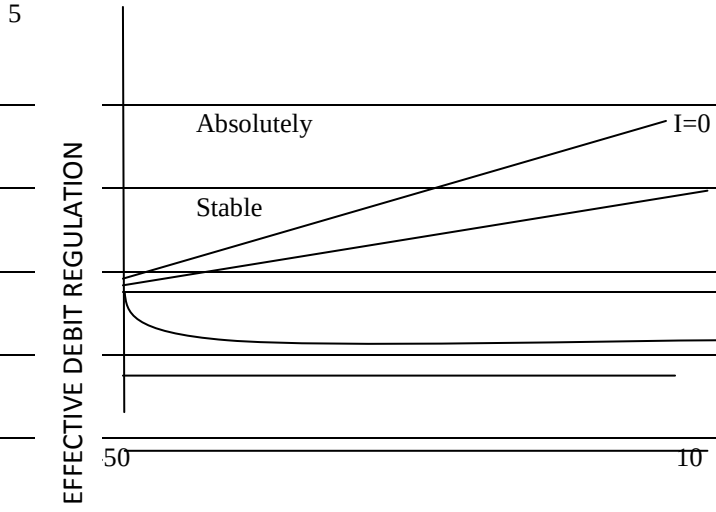



**DELAYED DEBIT EFFECT**

FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

## Utility Analysis And neural Network

Caroline Williams investigated affective density of the political and its effect on our understanding of political subjectivity. Taking up Spinoza's challenge to think about affect beyond corporeal embodiment, he argues that there is a modality of affectivity that cannot simply be inscribed within the borders of subjectivity. theorising affect as an impersonal force anchored in a relational ontology that gives due recognition to the circulation of affects, as well as to their ambivalent structure in creating sites of identification, and he utilises this ontology to reflect on the dynamic of the political and the shape of political subjectivity. Williams argues that Spinoza's philosophy (through ideas of conatus and imagination) offers the conceptual resources to reconfigure the composition of affective subjectivity as a transindividual social bond and as an unconscious dynamic of ethico-political existence. **Subjectivity (2010) 3, 245–262. doi:10.1057/sub.2010.15 Affective processes without a subject: Rethinking the relation between subjectivity and affect with Spinoza Caroline Williams** Phenomenological Time bears ample evidence and infallible observatory and apodictic evidence albeit with a sense of chagrined circumvention and defeated disgruntlment circularity of both psychological and phenomenological interpretations of art which are sometimes insightful, but most often work on the logic of cyclical time and unavoidable drives and representations played out as a ‘natural’ order: where ‘men can’t avoid being boys’ a circularity which valorizes a helplessness/truthfulness conundrum. One begins with a classical phenomenological fallacy with Duchamp’s Fountain, does the urinal keep returning or do we keep returning to the urinal? Put in this way, it is a male question, something that phenomenology conveniently conceals when positing the “I”, or indeed, the “we”. But we must be careful not to simply assume that the phenomenological first person is white, heterosexual, able-bodied and male, providing a foundation upon which knowledge should be based as a universal condition. Deleuze’s second critique of time is linked to Kant. This model of time extends not only to phenomenology but to the phenomenological interpretation

of art and is seen as a straight line. In the Critique of Pure Reason, Kant cuts up circular time and reassembles it as a series of sensory experiences which memory processes and which constitute subjectivity. Whereas the circular model of time (and interpretation) inscribes fatalism into the subject and the world, the phenomenological method is in danger of solipsism, where the subject constitutes the world. Much of embodied philosophy and enactivist theories of cognition go down this line with a plethora of art historical analyses by the likes of Rosalind Krauss, Michael Fried and others taking the lead from Merleau-Ponty. These approaches successfully challenge the **homunculus model of a Cartesian mind-body** dualism where there is a supervisor in the mind who processes experience and even consciousness and to whom all representations are directed. But phenomenology is mainly concerned with perceptual consciousness, the processing of colors, shapes, forms, motion in art and the world. The senses and sensorimotor processes are the fundamental pathways by which the interior is connected to the exterior in a chiasm of self-constitution in the world. The problem with this is that we explore the art through our sense perceptions, and what does art reveal to us?—our sense perceptions! Such solipsistic approaches have real problems relating complex and abstract conceptual production to bodily and sensorimotor contingencies, and they nearly always posit the subject as the site and object of discovery in the world. To be fair, phenomenology has moved on to rather more sophisticated models of intersubjectivity, taking into account the specificities of different kinds of bodies and ways of being beyond the white, male heterosexual as the common denominator, yet many naïve approaches to art perpetuate the notion that perceptual, rather than conceptual processes and evolutionary neural patterns are primary, fundamental ways to understand the art experience In fact, these aspects are only part of art production and reception, aspects which are often filtered out or made known to us in order for us to parody such trigger responses, as in much of Duchamp's work, which holds in contempt such retinal approaches to art. If we take the **subject of abstract expressionism** in the 1950s or later minimalist art, art history, teaching and interpretation still insists on putting the act of perceiving as fundamental to the constitution not only of the phenomenological first

person but also of art itself. The male act of usually standing up, which consists of directing a stream with the hand(s) into a container or on a flat surface, sand or snow in order to leave patterns and traces (the usual child's prank) underlies the logic of macho 1950s Abstract Expressionist culture, especially that of Jackson Pollock, so effectively parodied by Carolee Schneeman's Vagina Scroll, 1975. Here, painting or writing is a kind of a discharge which invests the **object and space with** the phenomenological "I", disciplined by existential struggle and reflected back by the traces, signatures and gestures left behind in the work of art, rediscoverable as an element in the making of other subjectivities.It is precisely this kind of solipsism which is blind to the art event as something well in excess of re-establishing habits of subjective self-identity through the production and interpretation of art. The attempt to reinforce a phenomenological identity upon the event by personalizing sense experience is unable to see art as something which remodels the phenomenological 'I' and neural plasticity itself. We learn new technologies and redraw our neural pathways through such new and spontaneous experiences which draw us into intersubjective worlds. Duchamp's Fountain and Deleuze's Repetition and Différence Grégory Minis sale (I have made some changes) The relation of different groups to deconstruction is settled by its relation to groups. An American enthusiast--one who has had a role in the development of our affection for Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^{\infty} \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^{\infty} cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$then given the constraints of $f(T_8)$, $(A1.5) can\ never\ be\ satisfied\ Co$nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8 \lambda_8 e^{-T_8 \lambda_8} + [(U_8 - V_8) M_8 + Y_8] - [((U_8 - V_8) M_8 -] e^{-T_8 \lambda_8} = 0$$

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$\boldsymbol{Z}_8^2 + M_8 Z_8 \sin(Z_8 T_8) - (A_8 M_8 + Y_8) + (A_8 M_8 + X_8)\cos(Z_8 T_8) = 0$$

$$(A_8 M_8) Z_8 - M_8 Z_8 \cos(Z_8 T_8) + (A_8 M_8 - X_8)\sin(Z_8 T_8) = 0$$

Where $A_8 = U_8 - V_8$Using (A2.3), substitute for $\boldsymbol{Z}_8^2$ and $Z_8$in (A2.2):

$W_8^2 (X_8^2 - 3 M_8 A_8 X_8 - M_8^2 (X_8 - Y_8) + 2 M_8^3 A_8 + 2 A_8^2 M_8^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_8 = 1 - \cos(Z_8 T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8 / 2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_8 \le 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2 (X_8 - Y_8)$$

$$Q_8(2, A_8) = (A_8 + 2M_8)^2 (2A_8 M_8 - (X_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2\left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8\left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8\left(2X_8 + Y_8\right)\right)$$

$$+2A_8\left(X_8 + Y_8\right)$$

$$=2M_8^2 W_8^2 (M_8 + A_8) + M_8 A_8^2 W_8 + (2A_8 M_8 - X_8 + Y_8)(A_8 W_8 + M_8 W_8^2$$

$$+ (2 - W_8)((2M_8 X_8 + M_8 Y_8) W_8 + A_8 (X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $W_8$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in *A* will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range $0 \le W_8 \le 2$.Defining

$Q_8(W_8, A_8) = a_8 W_8^2 + c_8$ *where* $a_8$ $b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \square - b_8 \square 0$$

And

$$b_8^2 \square 4a_8 c_8$$

$$Case(i): Y_8 = 0$$

$Let\ Y_8 = 0$ And try the substitution A= $X_8 = /2M_8\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8 /$ $2M_8$) is positive for *W<2. Since* $Q_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\ X_8$ ,

$Let\ Y_8 = \ y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 + \ n_8 y_8)$ we have

$$W_8^2\ y_8\left[X_8 n_8 \frac{(1 + n_8 y_8)}{(2 + n_8 y_8)} + M_8^2\left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8 X_8 \left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2}+2y_8\left(n_8-1\right)+\frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)}\right]=0.\ (A2.12)$$

Since $n_8$ and $A_8$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8=(2+n_8 y_8)$, *the condition*$(A2.88?)$*can be written as*

$$\left[\frac{X_8(E_8-8)}{E_8^2}\right]^2\left[\frac{X_8^2}{M_8^2}-4y_8 E_8^2\frac{X_8}{M_8^2}+4y_8(n_8-2)E_8^3\right]$$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8/W_8^2)$ can be expressed as

$$2E_8\left[y_8E_8+(y_8E_8(8+y_8)(E_8-n_8))^{8/2}\right]>\frac{X_8}{M_8^2}>2E_8\left[y_8E_8-y_8E_8(8+y_8)(E_8-n_8))^{8/2}\right]$$

Thoughts on the Social and Political Implications of Correlationism .Lion Mirror Truth be told, as my thought has evolved the issue of correlationism had fallen off the radar for me. Somehow the debate had come to seem too "philosophical" to me, too "scholastic", too remote from what interests me: understanding why social assemblages are organized as they are, how power functions in social assemblages, and what we might do to address that power and change things. Somehow the question of whether or not we can get out of the correlation between thinking and being just came to seem remote from these sorts of issues. Somehow it seemed too epistemological. ASIDE: Numerous discussions over the years have led me to believe that the debate over correlationism is poorly understood (or maybe I just don't understand it). On countless occasions I've heard people say "of course we must relate to things in order to know them." Well yeah, of course! I don't think this is what the critic of correlationism is getting at. It seems to me that correlationism is something more robust than the theses that we must relate to something to know it.

Correlationism instead seems to require the theses that thought and being are indiscernible. Put more concretely, the correlationist is someone who argues that we either a) can never tell whether being is merely a construction of our thought (weak correlationism), or b) who argues that thought actually constructs being (strong correlationism). In other words, correlationism is another name for idealism. One can hold that we must relate to something in order to know it without being a correlationist. As an aside I should also add that I am a correlationist about some things. For example, I think money is something constructed by society and is therefore a strong correlationist when it comes to money. At any rate, for a long time I'd become rather indifferent to debates about correlationism and philosophies of access. I had learned the lessons of speculative realism– which I could have also learned, I think, from Deleuze and Guattari, the new materialist feminists, actor-network theorists such as Latour –and had moved on. However, occasionally you come across a tone of phrase that pitches something in a different light. In The Cut of the Real, Katerina Kolozova writes, …the political problem of contemporary philosophy identified by the 'new realists' is, in fact, the product of a more fundamental epistemic problem. In his book After Finitude, Quentin Meillassoux calls this problem 'correlationism' and identifies it as an essentially post-Kantian legacy, which continues to dominate and limit philosophy. As a matter of fact, correlationism lies at (eb) the heart of postmodern theory and consists in the premise that thought can only 'think itself,' that the real is inaccessible to knowledge and human subjectivity, and that there is nothing but discursive constructs that fully determine thinking and that are meth0dologically accounted for all the way down. (1 – 2) Thought thinking only itself: Thought only encountering itself. In the jargon of postmodern and poststructuralist lingo, this would be the thesis of infinite semiosis, where signs ("thoughts") only ever relate to other signs. Within this framework, discursivity comes to be the hegemonic framework defining all of being. At the level of politics and social theory more generally, if the correlationist thesis is true the

consequences are clear: all social phenomena are discursive and all solutions to social and political problems will be discursive. The sole sphere of the political will be the discursive and all questions of politics will be questions of speech-acts and interpretation. The problem here is not that many theorists recognize that the discursive and semiotic plays an important role in the social and the political. It does and I've repeated this tirelessly. The problem is with what happens when thought or the semiotic becomes a hegemon, an "all", foreclosing our ability to recognize other forms of power. What I've wanted to say is that not all power functions discursively. In my last post and elsewhere I spoke of some other forms of politics: Thermopolitics: The politics surrounding **energy in the form of calories and fuels such as gasoline** and coal, and how our life and our very bodies are structured by energy dependencies and by being trapped in particular distributive networks that render these forms of energy available. I'm being quite literal when I speak of energy, talking about the effects, for example, of the absence of food in certain educational environments on cognition, for example; and am generally hostile to metaphorical extensions of the concept of energy which I see as erasing the dimension of real materiality. Geopolitics: The role that features of natural and built geography such as mountain ranges, rivers, oceans, soil conditions, roads, housing design, etc., play in the form that social relations take and how they impact individual bodies. Chronopolitics: The way in which the structurations of time organize what is possible for us. For example, the structuration of the working day, how much we can say and comprehend at any given time, the impact of things like the invention of the clock, etc. **Oikopolitics: This would be the domain of political economy described so well by Marxists**. So five different types of politics: Semiopolitics (or what currently dominates critical theory), thermopolitics, geopolitics, chronopolitics, and oikopolitics. No doubt there are other sites of the political or political struggle that we could speak of, but this is a good start. Also, it should be obvious that these aren't exclusive domains, but are entangled in all sorts of important

ways. For example, something might take place at the level of semiopolitics (speech, law, rhetoric, norms, and communication) that has all sorts of impact at the level of thermopolitics. Congress might decide to cut programs that fund school meal programs. This, in turn, will have a thermodynamic impact on those students that go without the calories they need developmentally and cognitively to function in a particular way. There is an entanglement here of semiopolitical and thermopolitical domains. The young student here has been constrained both at the level of **semiotic phenomena and thermodynamic structures**. The point is that if true, semiotic intervention (speech-acts, protests, interpretations, deconstructions, etc) will not be an appropriate response to all political problems because social formations are not entirely structured by the semiotic. The child in that school does not suffer from a lack of the right signs, but from a lack of calories needed to run the engine of his thought and body. Certainly semiotic interventions might be needed to render that energy available, but it is the energy itself that is at issue and the absence of that energy that forms the spider web entangling him in his position. A correlationist perspective tends to erase this as even being a site of the political. **Larval Subjects January 28, 2014 Ontological Anarché Posted by larval subjects under Uncategorized January 27, 2014 Roden on Pluralism Over at enemy industry Roden has an excellent post up on the pluralism discussion. Check it out here. January 25, 2014 Different Senses of Pluralism and Ontology Posted by larval subjects under Uncategorized Responding to one of my comments over at his blog, Jeremy Trombley:** I have deleted restricted, reformulated some sentences .Kindly bear with me. Intention is to make the discursive enucleation consistent with the subject in question**.** Recourse to epistemology of the quantum mechanics generic and fervor of string theory is a sine qua non for the further comprehension of the theory and ontology .How would the world appear to us if its ontology was that of classical mechanics but every agent faced a restriction on how much they could come to know about the classical state? We show that in most respects

it would appear to us as quantum. The statistical theory of classical mechanics, which specifies how probability distributions over **phase space evolve** under Hamiltonian evolution and under measurements, is typically called Liouville mechanics, so the theory we explore here is Liouville mechanics with an epistemic restriction. The particular epistemic restriction we posit as our foundational postulate specifies two constraints. The first constraint is a classical analog of Heisenberg's uncertainty principle; the second-order moments of position and momentum defined by the phase-space distribution that characterizes an agent's knowledge are required to satisfy the same constraints as are satisfied by the moments of position and momentum observables for a quantum state. The second constraint is that the distribution should have maximal entropy for the given moments. Starting from this postulate, authors derive the allowed preparations, measurements, and transformations and demonstrate that they are isomorphic to those allowed in Gaussian quantum mechanics and generate the same experimental statistics. They argue that this reconstruction of Gaussian quantum mechanics constitutes additional evidence in favor of a research program wherein quantum states are interpreted as states of incomplete knowledge and that the phenomena that do not arise in Gaussian quantum mechanics provide the best clues for how one might reconstruct the full quantum theory. I: http://dx.doi.org/10.1103/PhysRevA.86.012103 **Reconstruction of Gaussian quantum mechanics from Liouville mechanics with an epistemic restriction Phys. Rev. A 86, 012103 – Published 10 July 2012 Stephen D. Bartlett, Terry Rudolph, and Robert W. Spekkens.** Deleuze in fact characterizes modes of existence, with their powers and capacities. The answer is this: Deleuze approaches modes of existence, ethically speaking, not in terms of their will, or their conscious decision making power (as in Kant), nor in terms of their interests (as in Marx, for example), but rather in terms of their drives. For Deleuze, conscious will and preconscious interest are both subsequent to our unconscious drives, and it is at the level of the drives that we have to aim our ethical analysis. Here, I would like to

focus on two sets of texts on the drives taken, not from Nietzsche and Spinoza, but rather from Nietzsche and Leibniz (Leibniz being one of the first philosophers in the history of philosophy to have developed a theory of the unconscious). The first set of texts comes from Nietzsche's great early book entitled Daybreak, published in July 1881. Nietzsche first approaches the question of the drives by giving us an everyday scenario: "Suppose we were in the market place one day," he writes, "and we noticed someone laughing at us as we went by: this event will signify this or that to us according to whether this or that drive happens at that moment to be at its height in us—and it will be a quite different event according to the kind of person we are. One person will absorb it like a drop of rain, another will shake it from him like an insect, another will try to pick a quarrel, another will examine his clothing to see if there is anything about it that might give rise to laughter, another will be led to reflect on the nature of laughter as such, another will be glad to have involuntarily augmented the amount of cheerfulness and sunshine in the world—and in each case, a drive has gratified itself, whether it be the drive to annoyance, or to combativeness or to reflection or to benevolence. This drive seized the event as its prey. Why precisely this one? Because, thirsty and hungry, it was lying in wait" (D 119). This is the source of Nietzsche's doctrine of perspectivism ("there are no facts, only interpretations"), but what is often overlooked is that, for Nietzsche, it is our drives that interpret the world, that are perspectival—and not our egos, not our conscious opinions. It is not so much that I have a different perspective on the world than you; it is rather that each of us has multiple perspectives on the world because of the multiplicity of our drives—drives that are often contradictory among themselves. "Within ourselves," Nietzsche writes, "we can be egoistic or altruistic, hard-hearted, magnanimous, just, lenient, insincere, can cause pain or give pleasure" (Parkes, pp. 291-292). We all contain such "a vast confusion of contradictory drives" (WP 259) that we are, as Nietzsche liked to say, multiplicities, and not unities. Moreover, these drives are in a constant struggle or combat with each other:

my drive to smoke and get my nicotine rush is in combat with (but also coexistent with) my drive to quit. This is where Nietzsche first developed his concept of the will to power—at the level of the drives. "Every drive is a kind of lust to rule," he writes, "each one has its perspective that it would like to compel all the other drives to accept as a norm" (WP 481). DELEUZE AND THE QUESTION OF DESIRE: TOWARD AN IMMANENT THEORY OF ETHICS Daniel W. Smith. **In on a Possible Physical Metatheory of Consciousness Miroljub Dugic et al** show that the modern quantum mechanics, and particularly the theory of decoherence, allows formulating a sort of a physical metatheory of consciousness. Particularly, the analysis of the necessary conditions for the occurrence of decoherence, along with the hypothesis that consciousness bears (more-or-less) well definable physical origin, leads to a wider physical picture naturally involving consciousness. This can be considered as a sort of a psycho-physical parallelism, but on very wide scales bearing some cosmological relevance. Open Systems & Information Dynamics Volume 9, Number 2, 153 (2002) arXiv: quant-ph/0212128. On a broader perspective and background, the theoretical computation of the universe is justified.IN physics and cosmology, digital physics is a collection of theoretical perspectives based on the premise that the universe is, at heart, describable by information, and is therefore computable. Therefore, according to this theory, the universe can be conceived of as either the output of a deterministic or probabilistic computer program, a vast, digital computation device, or mathematically isomorphic to such a device. Digital physics is grounded in one or more of the following hypotheses; listed in order of decreasing strength. The universe or reality: is essentially informational (although not every informational ontology needs to be digital); is essentially computable (the pancomputationalist position); can be described digitally; is in essence digital; is itself a computer (pancomputationalism); is the output of a simulated reality exercise. Following Jaynes and Weizsäcker, the physicist John Archibald Wheeler wrote the following: [...] it is not unreasonable to imagine that

information sits at the core of physics, just as it sits at the core of a computer. (John Archibald Wheeler 1998: 340) It from bit. Otherwise put, every 'it'—every particle, every field of force, even the space-time continuum itself—derives its function, its meaning, its very existence entirely—even if in some contexts indirectly—from the apparatus-elicited answers to yes-or-no questions, binary choices, bits. 'It from bit' symbolizes the idea that every item of the physical world has at bottom—a very deep bottom, in most instances—an immaterial source and explanation; that which we call reality arises in the last analysis from the posing of yes–no questions and the registering of equipment-evoked responses; in short, that all things physical are information-theoretic in origin and that this is a participatory universe. (John Archibald Wheeler 1990: 5) David Chalmers of the Australian National University summarised Wheeler's views as follows: Wheeler (1990) has suggested that information is fundamental to the physics of the universe. According to this 'it from bit' doctrine, the laws of physics can be cast in terms of information, postulating different states that give rise to different effects without actually saying what those states are. It is only their position in an information space that counts. If so, then information is a natural candidate to also play a role in a fundamental theory of consciousness. We are led to a conception of the world on which information is truly fundamental, and on which it has two basic aspects, corresponding to the physical and the phenomenal features of the world. (Wikipedia). **Participatory consciousness forms the bastion, pillar, post, stylobate and fulcrum of the redressal measures and corrective steps taken for circumvention of many problems, identified and unidentified, spread widest commonalty. Anne Baring's sententious and pithy article will be good to have a timely outlook.** I would like to open my contribution to this meeting with the words of Ken Wilber that I quoted in the first Millennium Symposium: 'The secret impulse of life is towards greater consciousness. Maybe the evolutionary sequence is from matter to body, to mind, to soul, to spirit - each transcending and including, each with a greater depth and greater

consciousness and wider embrace' Ken Wilber, and with these from Rupert Sheldrake's contribution to the same meeting: "It is possible to make quite a good case that the sun could be conscious and if the sun, then why not the stars? Why not the galaxies? Why should there not be galactic minds? And if they communicate with each other then the distances are far too great for light to be the medium of transmission. We are into inter-galactic telepathy." After many decades of study, I know of only one diagram in all religious traditions that can adequately explain these statements; that can give us both a blue-print of the many levels of consciousness and a route-map to help us reach beyond our present level of understanding. I would like to introduce a diagram of the four worlds of Kabbalah to illustrate the theme that spirit has brought these four worlds into being and that we, living in the fourth and manifest world know nothing of the existence of the other three invisible ones. It is the clearest diagram I know of that can explain to us the multi-levelled structure of consciousness and our place in the great chain of being. We are indeed, as Wilbur suggests, evolving from matter to body, to mind, to soul, to spirit but where have we come from? It is possible that, as participants in the emanation of spirit into manifest life, we have also come from spirit to soul, to mind, to body, to matter. Everything, from the immense galactic energies to the minute forms of matter participates in one life, is generated from one creative spirit. There is nothing outside this life, this spirit. This diagram suggests that truth is not an article of faith but the experience of participating in the undiscovered reality of what has brought us into being. We cannot know truth until we have entered into communion with that reality. Continuing in similar vein **Michel Bitbol dilates upon the dialectic deliberation and polemical argumentation between** hermeneutists and eliminativists. When he formulated the program of Neurophenomenology, Francisco Varela suggested a balanced methodological dissolution of the “hard problem” of consciousness. He shows that his dissolution is a paradigm which imposes itself onto seemingly opposite views, including materialist approaches. I also point out that

Varela's revolutionary epistemological ideas are gaining wider acceptance as a side effect of a recent controversy between hermeneutists and eliminativists. Finally, he emphasizes a structural parallel between the science of consciousness and the distinctive features of quantum mechanics. This parallel, together with the former convergences, point towards the common origin of the main puzzles of both quantum mechanics and the philosophy of mind: neglect of the constitutive blindspot of objective knowledge. Phenomenology and the Cognitive Sciences Phenomenology and the Cognitive Sciences **Phenomenology and the Cognitive Sciences 2002, Volume 1, Issue 2, pp 181-224 Science as if situation mattered Michel Bitbol**

## ABSOLUTE STABILITY

Hence if $X_8 / M_8^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_8$ ) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8 - 8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2 - 2y_8\frac{(E_8 + n_8^2 y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right) - 4y_8(E_8 - n_8)$$

And has the solution

$$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8 - 8)}[y_8(E_8 + n_8^2 y_8) + [y_8^2 (E_8 + n_8^2 y_8)^2 + 4 y_8 (E_8 - 8)(E_8 - n_8)]^{8/2}]$$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2E_8[y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSOILIDATION**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + (M_8 + U_8)\lambda_8 - (M_8 + V_8)\lambda_8 e^{2T_8\lambda_8} + (M_8 U_8 + Y_8)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8\, e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + (M_8 + V_8) Z_8 \sin(Z_8 T_8) - (U_8 M_8 + Y_8) + (M_8(U_8 + V_8) - X_8)$$

$$\mathrm{X}\cos(Z_8 T_8) - 2M_8 V_8 \cos^2(Z_8 T_8) + M_8 V_8 = 0$$

$$(M_8 + U_8) Z_8 - (M_8 + V_8) Z_8 \cos(Z_8 T_8) + (M_8(U_8 + V_8) - X)$$

$$\mathrm{X}\sin((Z_8 T_8) - M_8 V_8 \sin(Z_8 T_8) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 =$

$1 - \cos Z_8\, T_8$and $A_8 = U_8 - V_8$we obtain

$$4M_8^2 V_8^2 W_8^4 + 2M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8(M_8 + V_8) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2[M_8 - V_8)^2(2A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8(M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8(A_8^2 + 4(X_8 + Y_8))]$$

$$+ W_8 \left[A_8^2 + 2X_8\right)\left(2A_8 M_8 - X_8\right) + 2A_8 Y_8 (M_8 + V_8) - M_8 A_8^3] + A_8^2 \left(X_8 M_8 - Y_8\right) = 0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8 > 0$$

$$and$$

$$C_8 > (X_8 - Y_8)/2M_8$$

$recalling\ that\ A_8 = U_8$ and defining $C_8 = U_8 + V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8(0, U_8 V_8 = A_8^2 (X_8 + Y_8)$$

$$Q_8\left(2, U_8 V_8\right) = \left(C_8 + 2M_8^2\right)^2 (2C_8 M_8 + Y_8 - X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$the existence of a D-position at this point is determined by letting cos $(Z_8 T_8) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all $T$, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**ABSOLUTE STABILITY IN PREDATOR- PREY MODELS**

**APPLIED TO ACCETUATION ATTRITION MODELS**

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability

criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The Percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending

upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t-T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\big(R(t)\big)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The

regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the

**TIME DELAY**

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a dominant eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

DELAY DUE TO CONGREGATIONAL AND CONFEDERATIONAL DEBIT CREDIT TYEPES

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t-T_8)F_8(R_8(t-T_8) - N_8(t)M_8((R_8(t) \quad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8} + \frac{N_8 dG_8}{dR_8} - \frac{dB_8}{dR_8} > \frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8} + \frac{dM_8}{dR_8}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_8 = \frac{N_8 G_8}{2M_8}\frac{d(8nF_8)}{dR_8} (\geq 0) \quad (8a)$$

$$I_8 = -\frac{N_8 G_8}{2M_8}\frac{d(8nM_8}{dR_8} (\geq 0) \quad (8b)$$

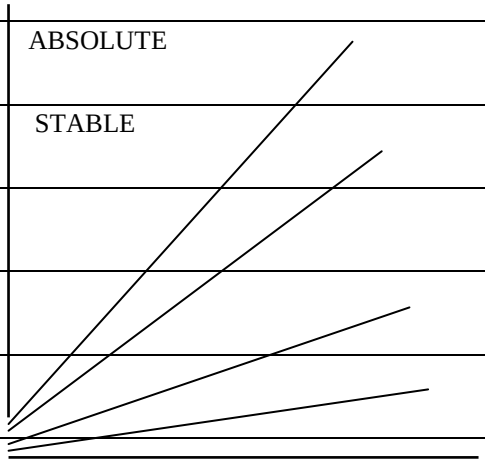
 

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8 > L_8 - I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_8 > L_8 - I_8, \quad \text{when} \quad P_8^2/4 > I \tag{8a}$$

$$P_8 > L_8 - I_8 + \frac{(I_8 - P_8^2}{(L_8 - I)} \quad \text{when} \quad P_8^2/4 < I. \tag{8b}$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit

(like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

See the diagram Nunney gives for predator and prey analysis. We can simulate the same for the financial system in the foregoing too.

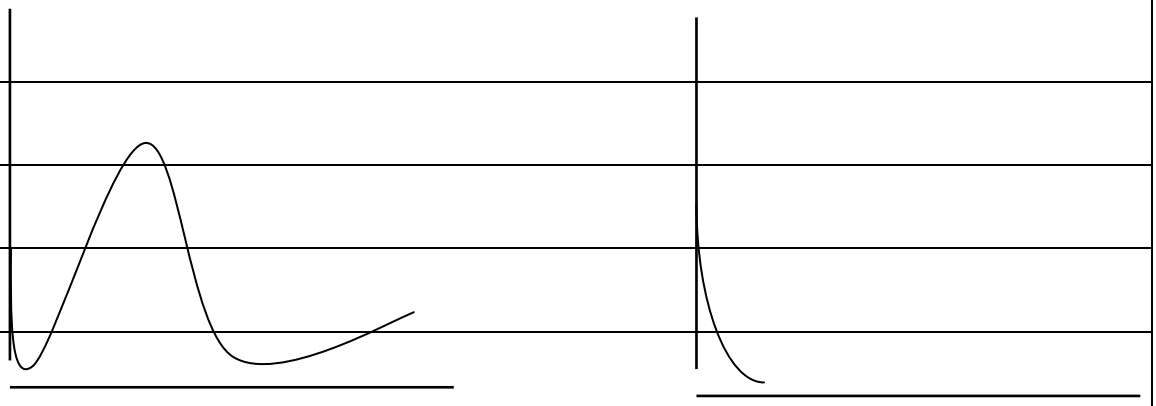
Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the

debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

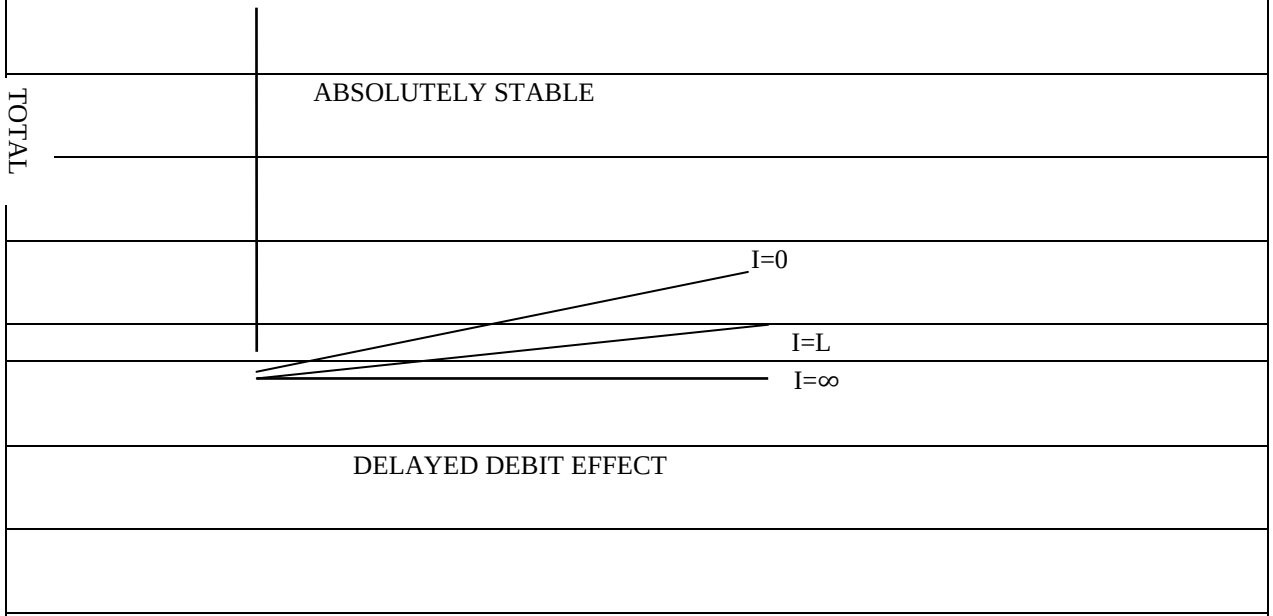
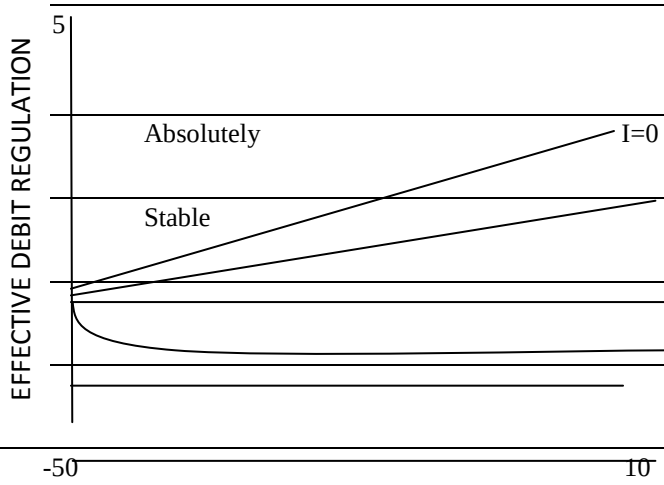



FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison?

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are

defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^\infty \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$ then given the constraints of $f(T_8)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8 \lambda_8 e^{-T_8 \lambda_8} + [(U_8 - V_8) M_8 + Y_8] - [((U_8 -$$

$$V_8 M_8 -] e^{-T_8 \lambda_8} = 0$$

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$\boldsymbol{Z}_8^2 + M_8 Z_8 \sin(Z_8 T_8) - (A_8 M_8 + Y_8) + (A_8 M_8 + X_8)\cos(Z_8 T_8) = 0$$

$$(A_8 M_8) Z_8 - M_8 Z_8 \cos(Z_8 T_8) + (A_8 M_8 - X_8)\sin(Z_8 T_8) = 0$$

Where $A_8 = U_8 - V_8$Using (A2.3), substitute for $\boldsymbol{Z}_8^2$ and $Z_8$in (A2.2):

$W_8^2 (X_8^2 - 3 M_8 A_8 X_8 - M_8^2 (X_8 - Y_8) + 2 M_8^3 A_8 + 2 A_8^2 M_8^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_8 = 1 - \cos(Z_8 T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8 / 2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_8 \le 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2 (X_8 - Y_8)$$

$$Q_8(2, A_8) = (A_8 + 2M_8)^2 (2A_8 M_8 - (X_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2\left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8\left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8\left(2X_8 + Y_8\right)\right)$$

$$+2A_8\left(X_8 + Y_8\right)$$

$$=2M_8^2 W_8^2 (M_8 + A_8) + M_8 A_8^2 W_8 + (2A_8 M_8 - X_8 + Y_8)(A_8 W_8 + M_8 W_8^2$$

$$+ (2 - W_8)((2M_8 X_8 + M_8 Y_8) W_8 + A_8 (X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $W_8$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in *A* will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range $0 \le W_8 \le 2$.Defining

$Q_8(W_8, A_8) = a_8 W_8^2 + c_8$ *where* $a_8$ $b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \square b_8 \square 0$$

And

$$b_8^2 \square 4a_8c_8$$

$$Case(i): Y_8 = 0$$

$Let\ Y_8 = 0$ And try the substitution A= $X_8 = /2M_8 in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8 /$ $2M_8$) is positive for *W<2. Since* $Q_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\ X_8$ ,

$Let\ Y_8 = \ y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 + \ n_8 y_8)$ we have

$$W_8^2\, y_8\left[X_8 n_8 \frac{(1 + n_8 y_8)}{(2 + n_8 y_8)} + M_8^2\left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8 X_8 \left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2}+2y_8\left(n_8-1\right)+\frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)}=0.\ (A2.12)\right.$$

Since $n_8$ and $A_8$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8=(2+n_8 y_8),$ *the condition*(A2.88?)*can be written as*

$[\frac{X_8(E_8-8)}{E_8^2}]^2[\frac{X_8^2}{M_8^2}-4y_8 E_8^2\frac{X_8}{M_8^2}+4y_8(n_8-2)E_8^3]$ ?0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8/W_8^2)$ can be expressed as

$2E_8[y_8E_8+(y_8 E_8(8+y_8)(E_8-n_8))^{8/2}]>\frac{X_8}{M_8^2}>2E_8[y_8E_8-y_8E_8(8+y_8)(E_8-n_8))^{8/2}]$

**ABSOLUTE STABILITY**

Hence if $X_8/M_8^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0?$W_8$ ?2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8-8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2-2y_8\frac{(E_8+n_8^2y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right)-4y_8(E_8-n_8)\ ?0$$

And has the solution

$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8-8)}[y_8(E_8 + n_8^2 y_8) + [y_8^2 (E_8 + n_8^2 y_8)^2 + 4 y_8 (E_8 - 8)(E_8 - n_8)]^{8/2}]$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2 E_8 [y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + (M_8 + U_8)\lambda_8 - (M_8 + V_8)\lambda_8 e^{2T_8\lambda_8} + (M_8 U_8 + Y_8)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8\, e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + (M_8 + V_8) Z_8 \sin(Z_8 T_8) - (U_8 M_8 + Y_8) + (M_8(U_8 + V_8) - X_8)$$

$$\mathrm{X}\cos(Z_8 T_8) - 2M_8 V_8 \cos^2(Z_8 T_8) + M_8 V_8 = 0$$

$$(M_8 + U_8) Z_8 - (M_8 + V_8) Z_8 \cos(Z_8 T_8) + (M_8(U_8 + V_8) - X)$$

$$\mathrm{X}\sin((Z_8 T_8) - M_8 V_8 \sin(Z_8 T_8) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 = 1 - \cos Z_8 T_8$and $A_8 = U_8 - V_8$we obtain

$$4M_8^2 V_8^2 W_8^4 + 2M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8(M_8 + V_8) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2[M_8 - V_8)^2(2A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8(M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8(A_8^2 + 4(X_8 + Y_8))]$$

$$+ W_8[A_8^2 + 2X_8)(2A_8 M_8 - X_8) + 2A_8 Y_8(M_8 + V_8) - M_8 A_8^3] + A_8^2(X_8 M_8 - Y_8) = 0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8 > 0$$

$$and$$

$$C_8 > (X_8 - Y_8)/2M_8$$

$recalling\ that\ A_8 = U_8$ and defining $C_8 = U_8 + V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8(0, U_8 V_8 = A_8^2 (X_8 + Y_8)$$

$$Q_8\left(2, U_8 V_8\right) = \left(C_8 + 2M_8^2\right)^2 (2C_8 M_8 + Y_8 - X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$the existence of a D-position at this point is determined by letting cos $(Z_8 T_8) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all *T*, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**ABSOLUTE STABILITY IN PREDATOR- PREY MODELS**

**APPLIED TO ACCENTUATION ATTRITION MODELS**

We apply model to each and every module and a concatenated model was to be presented which we procrastinate due to spatial restrictions.

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - D(R) - NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

The question is political-discursive. It's what Stengers and Latour are getting at with their concepts of diplomacy and cosmopolitics. They grant, first, that all entities exist and, second, that to say that someone's cherished idol (or whatever disputed entity they hold dear) is non-existent is a **'declaration of war'** – 'this means war,' as Stengers often says. They thus shunt onto-political discourse off of the terrain of knowledge/belief in the sense of existence/non-existence. Their basic claim seems to be that 'respect for otherness,' i.e. political pluralism, can only come from granting the entities that others hold dear an ontology, even if you don't 'believe' in them. You

are thus permitted to say 'I do not follow that god, he has no hold over me' but you are not permitted to say 'your god is an inane, infantile, non-existent fantasy, grow up.' And it's not just a question of politeness (although there's that too). The point is to grant others' idols and deities an existence – one needn't agree over what that existence entails, over what capacities that entity has or what obligations it impresses upon you as someone in its partial presence but to deny it existence entirely is to 'declare war' – to deny the possibility of civil discourse, of pluralistic co-existence. I believe that it was Richard Rorty who once quipped something like claims to reconcile realism and idealism always seem to end with a triumph of idealism. We don't, in fact, get realism through such approaches, but rather just get a pervasive anti-realism. I think this is also the problem with the "non-controversial pluralism" advocated by Stengers and Latour that Phillip defends here. Such pluralism is not realism but is, in fact, a thoroughgoing social constructivism. I think this is the central problem with Latour's argument in Irreductions (these days I regret having ever defended it). In rejecting both Enlightenment critique and what he calls "reduction" he wants to say something like "The Pentacostal really is filled with (e&eb) the Holy Spirit", that for the 19th working scientist heat really is a fluid and phlogiston really is what allows things to burn, and that for the Greek lightning really is an expression of Zeus's anger. Latour tells us that we aren't to reduce or explain away the entities posited by another group's "ontology" but are to develop explanations from within that ontology. It might not sound particularly sexy– and it certainly doesn't tell us what is worth thinking –but I can't help but believe that philosophy is the critical and reflective investigation of basic concepts that guide our investigation of the world about us, how we ought to live our lives, and what form of governance might be best. Compare two figures. A scientist might ask, what causes depression? We can very well imagine a philosopher turning around and asking the scientist, what is causality? The scientist presupposes a concept of causality in her investigations. She uses this concept in her inquiry. Now she might have a sophisticated concept of causality or she might never have thought much about causality at all, using it in the sort of colloquial and unreflective way that Plato decried when, for example, people

like Euthyphro talked about piety. A whole cascade of questions arise when we raise a question like "what is causality?" We can ask whether or not causality exists at all. We can ask how we distinguish between correlation and two events that merely accompany one another from genuine causation. This, for example, was Hume's question. But perhaps most importantly we can ask whether there is only one form of causality or many forms of causality. Is there only one-to-one causation; one cause and one effect? Is there many-to-one causation; or many events conspiring to produce an effect? Is there one-to-many causation; or one event producing a variety of different effects? We can even ask whether causality necessarily moves from past to present or whether there aren't forms of causality that move from future to past! **January 23, 2014 Pluralism and Realism:** multiple-worlds Over at Struggles Forever, Jeremy Trombley has an interesting post up on "the ontological turn" in anthropology or ethnography. I've been meaning to have a discussion with him about this as I think it's an issue many of us are struggling with. For example, the core project of The Democracy of Objects– a project which I think many have missed –is to somehow reconcile some version of social constructivism with a realist ontology capable of making room for ecology (which requires realist and materialist positions as there's a fact of the matter where global warming is concerned) as well as the role played by objective agencies in social assemblages such as technologies, infrastructure, features of geography, local climates, the growth cycles of plants and animals, waste, etc. Maybe we can try to organize some cross-blog event to discuss these issues. I certainly think they're close to the heart of Jeremy, Michael of Archive Fire, Arran James, and a host of others. As an aside, I'm beginning to realize how the different sites of the political I've been outlining– semiopolitics, thermopolitics, oikopolitics (political economy), geopolitics, eropolitics (the politics of sex and desire), biopolitics, and chronopolitics (and I'm sure there are other political sites!) –are drawing me away from traditional Marxism. Assuming that classical Marxism holds that economics or the conditions and relations of production are determinative of all other sites of the political, the various sites of the political that I've been outlining would lead to the conclusion that there is not one determinative

base of the political. This would not require committing Marx to flames, but rather of recognizing the phenomenon of overdetermination, or of a variety of different entangled sites of the political. But I digress. First, I find myself wondering what the ontological turn means in ethnography. Is it 1) the investigation of the different ontologies held/proposed by different cultures? E.g., the Aztecs believed that reality was structured in this way, while the Greeks in that way, and the ancient Chinese this way, etc? Or 2) is it an investigation of how real entities– independent of cultural beliefs –influence cultural formations? Or is it a combination of both? A position that I would favor. January 20, 2014 Model is applied to structured reality

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

**The evaporation spectrum of black holes from a local quantum gravity perspective Aurélien Barrau (Submitted on 21 May 2016**)

(13) Authors revisit the hypothesis of (e) a possible line structure in the Hawking evaporation spectrum of black holes, due to (e) non-perturbative quantum gravity effects, even arbitrarily far away from (e) the Planck mass.

(14) Authors show that this naive prediction might in fact hold in (eb)the specific context of loop quantum gravity, with (e&eb) a small departure from (e) the ideal case for (e) some low-spin transitions.

They also show that the effect is (=) neither washed out by the dynamics of (e) the process, nor by (e) existence of a mass spectrum up to (e) a given width, nor by the secondary component induced by (e) the decay of neutral pions emitted during (e&eb) the time-integrated evaporation. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Cite as: arXiv:1605.06613 [gr-qc] (or arXiv:1605.06613v1 [gr-qc] for this version

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t - T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\left(R(t)\right)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

It's very easy to treat this as an abstract, academic debate: "Are you a realist or are you an anti-realist?", as if it were just a matter of what happens to be true. But it seems to me that this debate has, in reality, always been about politics. As I outlined in my last post, we have perpetually seen how appeals to the real and natural have been used in the name of oppressive power, inscribing both the exploitation of nature and the oppression of various people in the very fabric of being itself. Theistic theology and realist ontology have perpetually been used in the name of

what Deleuze called "State Philosophies" or philosophies that ontologize contingent orders of power and privilege (e.g., "the great chain of being" used to justify patriarchy, monarchy, serfdom, poverty, etc, and appeals to nature used to justify poverty and racial inequality (The Bell Curve), patriarchy (evolutionary psychology), heteronormativity, capitalism, etc). Because arrangements of power and inequality are always contingent in the sense that there's no marked difference in the capacities of peoples, power always looks for a transcendent supplement that would provide justification through ontological necessity. Antirealism– from the Greek atomists to present –became the radical and emancipatory gesture because it revealed the lie behind all of these forms of social organization or their inherit contingency or arbitrariness. Realism, by contrast, has all too often functioned as an apologetics for arbitrary power and social organizations. Here it's worth recalling what Foucault said about science: "…Even before we know to what extent something like Marxism or psychoanalysis is analogous to a scientific practice in its day-to-day operations, in its rules of construction, in the concepts it uses, we should be asking the question, asking ourselves about the aspiration to power that is inherent in the claim to being a science. The question or questions that have to be asked are: "What types of knowledge are you trying to disqualify when you say that you are a science? What speaking subject, what discursive subject, what subject of experience and knowledge are you trying to minorize when you say 'I speak this discourse, I am speaking a scientific discourse, and I am a scientist.' What theoretico-political vanguard are you trying to put on the throne in order to detach it from the massive, circulating, and discontinuous forms that knowledge can take?" (Society Must Be Defended, 10). All of these questions hold equally for claims to something being real. What is one trying to minorize when claiming something is real? What becomes privileged? What is excluded? It is these questions that have been at the heart of the realism debates, for as Spencer-Brown taught us, every distinction has a marked and unmarked space, draws attention to something to be included and pushes something into the unconscious or the domain of the invisible, hidden, or veiled. This is above all the case with evocations of the real. However, as I've tried to show antirealism

leads to its own problems. First, so long as we exclude real beings from our ontological inventory, we are unable to fully understand how power functions ("such and such a set of cultural formations have been thoroughly debunked, yet people still live as if they believed them"). Not only do we not fully understand the sources of the problems due to too much focus on the discursive and semiotic, but we deny ourselves valuable sites of political intervention at the level of infrastructure. Second, we are prevented from addressing things such as cultural racism, such as that found in Heidegger's privileging of the West and the Greeks and Germans in particular. We do a good job addressing biological and theological racism, heteronormativity, and sexism by showing how it is a cultural construction, but when faced with racism such as Heidegger's where he argues that there's something "unique" about the Greek event, it's language, and about the German language, or Badiou/Zizek's racism with respect to the "Pauline Event", we really have no response. Here someone like Jared Diamond or Fernand Braud el is needed to explain global-geographical inequalities. Third, the tools of the cultural turn really do not provide us with the means of thinking the ecological as a site of the political. For this reason, I've tried to formulate a third way– which might be called "constructivist realism" or "constructivist naturalism" –that retains the insights of the cultural turn, while also allowing a robust place for the material. I don't claim to be original in this. I think that many such as Manuel DeLanda, Deleuze and Guattari, Stacy Alaimo, Karen Barad, Jane Bennett, etc., are up to something similar. It's a vast project that requires the work of a multitude of voices, especially given the way in which the culturalism pervades humanities. **The Relational and the Non-Relational: Notes towards an Immanent and Pluralist Theory of Meaning Posted by larval subjects under uncategorized k-bigpic I have deleted or reformulated some sentences for which I beg pardon attributable to spatial constraints** What is correlation, what is causation? What is coordination? What is coincidence? These thoughts have bothered humans since long. I have been able to excavate some reliable material from Wikipedia which I would like to share. "All views are partial and contingent – that's the lesson of pluralism, and you say as much yourself when you say that we must be attentive to ways our own

knowledge might contain superstitions. But it's not discourse that determines whether one view works and another doesn't – it's a confrontation with **non-discursive (non-human)** agencies." I wonder if part of the issue here is that we understand pluralism differently. For me this doesn't sound like pluralism at all. Rather, it just sounds like our **epistemological condition**. We [hopefully] come to understand those portions of existence we question or investigate. The reason we come to investigate them is largely contingent. Finally, the accounts of these features of existence we give can be mistaken. All of this is perfectly consistent with a monism. I'm not sure why what he outlines above is the lesson of pluralism. It seems to me that every realist knows this. It seems to me that ontological pluralism is something quite different. Ontological pluralism is the thesis that there are many different worlds inhabited by many different entities. Thus, for example, you would have one world; say that of Lucretius, that's only inhabited by atoms and their combinations. You would have another world; say the world of the Mongolian shaman, that's inhabited by spirits that do all sorts of things. The ontological pluralist is saying that these spirits are, that they exist, that they're real. This is something quite different than merely saying that there are partial and contingent points of view on the world. It's this saying that these entities are rather than that some person or groups of people believe that they are that is the nub of the issue. Now, I think part of the issue here is that there's an ambiguity in the term ontology. Ontology can be one of two things. On the one hand, ontology is a group or persons set of beliefs as to what is. Here it's trivially true that there are pluralities of ontologies and the realist readily recognizes this. This is the whole reason there are debates over ontology. Mongolian shamans have their ontology, Europeans theirs, Christian fundamentalists theirs, materialists theirs, etc. When striving to understand and communicate with others it's vital to understand these ontologies because, as rhetoricians like Burke point out, our **beliefs about what is** are among the things that motivate our action. Despite having never seen bacteria I was my hands because I believe there are bacteria and viruses on door handles and whatnot and don't want to get sick. My belief about a particular thing existing is what motivates my action. And who knows, perhaps this belief is as

superstitious as the belief that the crops failed because God was displeased with my community. On the other hand, ontology is a theory **about what is.** It is making a claim that something exists. This is where the rubber hits the road. The ontological pluralist seems committed to the thesis that every group set of beliefs about what exists is sufficient for granting the existence of those entities. This is what I find objectionable in Latour. Obviously I'm not bothered by Latour's suggestion that we should take into account the role that nonhumans like speed bumps, rivers, microbes, etc., play in the form that social assemblages take. But this is not all that Latour claims. This is precisely where the philosopher might balk. Unlike the ethnographer, the philosopher is not interested in what people believe exists, but rather philosophers– at least of the realist variant –are trying to figure out what is. In other words, the realist philosopher begins with the premise that not all of these beliefs about what is are true. So for the philosopher, recognizing that Mongolian shaman's believe in the existence of shamans would only be the first step. The next step would consist in determining whether there's good reason for thinking such entities really do exist, i.e., whether there's good reason for believing these entities have mind and culture independent reality. Lest readers think that I'm just picking on the supernatural here, we can ask similar questions about strings, subatomic particles, galaxies, etc., etc., etc. I guess the question really comes down to what exactly we mean by ontological pluralism. When we talk about ontological pluralism are we defending the thesis that people have different theories of being? If so, then ontological pluralism is trivially true. If this is what is meant, then I certainly share Jeremy's view that it's valuable to understand the different world's people believe in. Certainly when I was practicing as an analyst I didn't get in ontological debates with my patients and it was necessary to understand their theory of being or their ontology to properly attend to them. Or, when we talk about ontological pluralism are we defending the thesis that all these ontological theories are true and refer to really existing entities? That's quite a different claim and is not one I would defend or endorse. Now someone might object that "in both The Democracy of Objects and Onto-Cartography you defend the thesis that there are multiple worlds." This is true. Because I hold that not everything

is related I'm led to the conclusion that there are diverse worlds. However, I also hold that however many worlds there might be, these worlds are nonetheless composed solely of material entities. Within the framework I propose I wouldn't suggest that there's one world where there are spirits and another world where there are souls and yet another where there are only material entities. My view is that there aren't spirits or souls in any of these worlds. One of the things I keep hearing in these discussions is that somehow the realist adopts a view from nowhere. I honestly don't understand this criticism. Investigation always occurs somewhere and requires all sorts of mediations involving technologies, experiments, etc. It's that labor of gathering evidence, conducting experiments, using technologies to observe the world, etc., that gradually gives us a body of data that allows us to say there's good reason to believe that such and such a thing exists and has these powers. Another charge seems to be that the realist refuses to recognize that their claims about the world are fallible. I find this charge particularly strange because it's precisely because the realist recognizes the difference between our theories of the world or what we say about the world and the world itself that fallibility is built into the core of his position. Realism doesn't mean one holds they have special access to the world, that they know all truths, or that they have the truth in hand, only that there are truths to be known and that we can be mistaken about things. Circling **Squares Philip** has a post responding to my quandaries about how to mesh realism and pluralism. He writes: Ontologically and metaphysically the idea of realist pluralism is no longer an issue. There are (appropriately) numerous variants but the basic idea that reality is itself pluralistic is well established. Captioned expatiation is about Everett worlds and string theory

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from

being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

The von Neumann formulation leads us to conclude that the lower realms of quantum reality form a virtually limitless array of potentials that are inherently incapable of realization without the observation of a knowing entity, which we identify with consciousness (Kafatos and Nadeau, 1990). We thus have a correspondence between quantum theory, the hierarchical structure implied by the Holographic Principle, and the higher orders of experience,

leading to the full expression of the Conscious Universe. David Bohm's theory of the implicate and explicate orders involves a holographic Principle that is fully consistent with the Holographic Principle discussed here (Bohm 1980; 1986; Bohm and Hiley, 1993). According to Bohm, there is an implicate order that represents the universal, holographic subtext of reality, and which unfolds in every moment to produce the explicate order that we all observe. Thomas Germinario (2004) has equated the implicate order with (eb) the unconscious process, and the explicate order with (eb) conscious process. He emphasizes the nature of dreams within (eb) the implicate order, and the importance of the dream work for maintaining a healthy mind through (e&eb) integration of the implicate order and our daily lives. Allan Combs and Mark Holland (1990) connected the implicate order or holomovement (Bohm, 1980) with (e&eb) Carl Jung is theory of synchronicity (Jung and Pauli, 1955), with (e&eb) the implicate order providing (eb) a holographic medium through which apparently disconnected individuals become connected. The principle of synchronicity, the instantaneous connection of people and events beyond (e) the senses, has been equated with the quantum-physical principle of non-locality (Combs and Holland, 1990; Germine, 1991), and has been proposed to be the fundamental mechanism of conscious process (Germine, 1991).Applying the Holographic Principle theory of mind, subjects may be connected synchronously in the manner of two individuals conversing through cellular phones, where the radio signals are transmitted through (e&eb) a distant satellite. The difference here is that the radio signals travel at the speed of light, and thus the there is some miniscule time lapse between the sending and receiving ends. The Holographic Principle implies (eb) a much more distant and instantaneous communication pathway, with (e&eb) the analogue of the satellite being the holographic boundary of the Universe, as well as a much richer communication, potentially involving a transmission of thoughts and feelings. Our ego-consciousness seems to mask (e) the universal relatedness implied by (e) the Holographic Principle, and it is perhaps only through transcendence of (e) the ego-consciousness that the higher orders of experience can become conscious. In the phenomenon of synchronicity, there seems to be a meaningful connection between (e&eb) individuals that

breaks through the barrier of ego-consciousness. Such a connection is reported by many individuals in the course of dreams, when (e) the ego-consciousness has been suspended, at times informing (e&eb) the dreamer of something that has happened in the life of a meaningfully-connected individual. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Impassioned-intense Connectionist neuroscience, in an effort to attain relevance to clinical practice, has recently given us a theory of the relationship of the unconscious to (e&eb) consciousness which is (=) dualistic (Viamontes and Beitman, 2007), an ad hoc theory of conclusions based on (e) assumptions, the most erroneous of which are that representations of (e) information are information, and that such representations elaborate themselves through (e&eb) a process of representations of representations in neuronal circuits (Bechtel and Abrahamsen,1991). As theory, the connectionist models are lacking (e) heuristic value and **are (=) inimical to the progress** of (e) basic and clinical neuroscience. This is not to denigrate or minimize the importance of neuronal connections, but these connections are rather like the wires in (eb) a radio that plays a symphony. The radio, in and of itself, is (=) incapable of playing (e&eb) the symphony. The radio doesn't do not write the symphony, nor does it broadcast the Passionate-raging symphony over the airways. Yet, if we loosen one wire, the symphony is (=) not manifested through the radio. The brain is (=) not subject to (e) the same vulnerability, in that it has (e) parallel and redundant networks of processing, and, it is in this sense that the connectivity of the brain becomes (eb) important. Connectionism, to the extent that it accepts (e) consciousness and the (e&eb) unconscious at all, attributes them to (e) separate groups of "circuits" in the brain, with (e&eb) consciousness-processing circuits having (e) a limited but detailed ability to analyze (e&eb) information, and unconscious-processing processing more information in less detail (Viamontes and Beitman, 2007). There is an element of truth in this view in that information that has importance and meaning resonates within the brain at higher levels of experience and at higher levels of holographic recursion. This resonance involves what we know, what we feel, and what we are capable and willing to consciously realize. The Universe, as William Blake noted,

can be seen in a grain of sand. Indeed, the grain of sand requires (e) the Universe to exist. The human organism requires (e) the Self to exist. The Self is not a simple amalgam of (e) sensory experiences, as suggested in the connectionist model (Blinder, 2007). This amalgam view of the Self can be damaging when applied to (e&eb) psychotherapy, as the connectionists suggest (Blinder, 2007). The Self is ours forever and ever, from eternity to eternity. As Milton suggested, the mind is (=) its own place. Our deepest purpose is for our minds to resonate with (e&eb) the supra-conscious levels of experience, much as it is (=) purpose of the radio to play the symphony. If the symphony is not to our liking, (eb) we can turn the radio off. So, perhaps, the transpersonal consciousness can turn (e&eb) our own individual and collective radios off, or perhaps turn the volume down, if (e) we are not in harmony with the **Universal Mind**. Physics has given us many enigmas. Quantum theory itself is an enigma. But some of the enigmas seem to involve (e&eb) the limitation of (e) our own abilities of discernment. The thought experiment of (e) "Schrödinger's Cat," as it is often interpreted (e.g. Gribbin, 1984), assumes (eb) that the **cat is not conscious, and is (=) incapable of observation in the context of (e) collapse of (e) the wave function.** This is due, in large part, to the false dichotomy of (e) information and (e&eb) experience. Our observation of other animals indicates, to those that are attuned to (e) the animal mind, who can "feel" the animal mind, that, at (eb) the very least, all mammals are (=) conscious. It would seem that the great neuropsychiatrist, Stanley Cobb (1948), may have been correct in his attribution of (e) levels of consciousness to a wide variety of organisms: "...lower animals with (e&eb) no cerebrum appear to be conscious...even plants such as the sunflower that turns towards strong light may have (e) a vague awareness or warmth and comfort. There are many degrees of consciousness and it is my contention that it is integrated at (e&eb) many levels like other important functions of (e) the central nervous system."With respect to the theory of mind, recent mainstream thinking remains (eb) classical and mechanistic. Such views of mind are advocated by our most in prominent neuroscientists (e.g. Changeux, 1985; Gazzaniga, 1985) as neuroscience research attempts to explain all mental function on the basis of (e) **pure mechanism and the localized function of**

**specific areas of the brain.** Francis Crick, the distinguished Nobel Laureate, had a very influential second career as a neuroscientist. He concluded that we are basically soulless creatures, that all mental processes could be reduced to (e&eb) identifiable, neural correlates in (eb) the brain, and that, in fact, we are (=) nothing but a collection of neurons (Crick, 1995). Crick even went so far as to say that what we call the soul is (=) a group of neurons located in the prefrontal lobes of the brain. Another Nobel Laureate, Gerald Edelman (1987), has written a critically acclaimed book on "neural Darwinism," which purports that the brain is "circuitry" is constructed through a developmental process of "survival of the fittest" neurons and neuronal connections. At the same time the total purposelessness and lack of direction in evolution is expounded, again to critical acclaim, by Richard Dawkins (1986). The idea that there essentially no self or soul imbued with agency and self determination was expounded by Gilbert Ryle (1949/2002) with his parody of the "ghost in the machine," and this work continues to be influential among many scholars of the mind. Philosophers such as Patricia Churchland (1986), who coined the term "neurophilosophy," take an "eliminative materialist" view of (e) belief, free will, and consciousness. Other philosophers, such as Daniel Dennett (1991) would make (eb) experience "epiphenomenal," a by-product of (e) brain processes with (e&eb) no effect whatsoever on (eb) the function of mind. However, unlike neuroscience, which is still in its infancy, the philosophy of mind has had a long history, including such notable figures as Plato, Aristotle, Descartes, Kant, Locke, Hume, Kierkegaard, and James, to name just a few. Perhaps, as noted by Ervin Laszlo (1974): "in today's world, most of the traditional functions of (e) cognitive synthesis have atrophied (e) and are ignored and neglected." Hopefully, what we have presented here will be a step in the right direction for what Laszlo calls a "conceptual synthesis" to help fill the need for (e) meaningful engagement is such a world. The fundamental elements of such a synthesis have been described by (e) Laszlo as follows: Conceptual synthesis performs (e&eb) at least five basic functions in the guidance of affairs. They are the mystical, the cosmological, the sociological, the pedagogical or psychological, and the editorial functions. The mystical function inspires (eb)

in man a sense of mystery and profound meaning related to (e&eb) the universe and of himself in it. The cosmological function forms (eb) images of the universe in accord with (e&eb) local knowledge and experience, enabling (eb) men to describe and identify (e&eb) the structure of (e) the universe and the forces of nature. The sociological function validates supports and enforces (eb) social order, representing (eb) it in accord with (e&eb) the nature of the universe, or as the natural or right form of (e) social organization. The pedagogical or psychological function guides (eb) individuals through stages of life, teaching ways of understanding themselves and others and presenting (eb) desirable responses to life is challenges and trials. Finally, the editorial function of (e) conceptual is to define (eb)some aspects of reality as important and credible and hence to be attended to, and other aspects unworthy of serious attention. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Kant sums up this inversion and its spirit early in the Critique of Pure Reason: Up to now it has been assumed that all our cognition must conform to the objects; but all attempts to find something about them a priori through concepts that would extend our cognition have, on this presupposition, come to nothing. Hence let us once try whether we do not get farther with the problems of metaphysics by assuming that the objects must conform to our cognition, which would agree better with the requested possibility of an a priori cognition of them, which is to establish something about objects before they are given to us. This would be just like the first thoughts of Copernicus, who, when he did not make good progress in the explanation of the celestial motions if he assumed that the entire celestial host revolves around the observer, tried to see if he might not have greater success if he made the observer revolve and left the stars at rest. (B xvi) Onticology, like all variations of object-oriented ontology, is realist in its orientation. In defending a realist ontology onticology holds that the vast majorities of objects, actants, beings, or entities are independent of humans and are what they are regardless of whether any humans regard them or register them. In short, onticology rejects any anthropomorphic, idealist, or anti-realist thesis to the effect that to be is to be the correlate of mind, spirit, the body, the human, and language or otherwise. While it is certainly the case that knowledge is necessarily dependent on the object to which it relates, the converse does not hold true. Objects are not dependent on being known, regarded, perceived, or spoken

about. As such, and to put it in Aristotlean terms, knowledge is an accident of objects, not objects an accident of knowledge. As Althusser so nicely puts it, "[n]o doubt there is a relation between thought-about-the-real and thisreal, but it is a relation of knowledge, a relation of adequacy or inadequacy of knowledge, not a real relation, meaning by this a relation inscribed in that real of which the thought is the (adequate or inadequate) knowledge" (Reading Capital, 96). Althusser goes on to remark that "[t]he distinction between a relation of knowledge and a relation of the real is a fundamental one: if we did not respect it we should fall irreversibly into either speculative or empiricist idealism" (ibid.). Onticology categorically endorses Althusser's verdict. It is a fundamental necessity to distinguish between those relations that belong to the object and those that belong to knowledge. Contemporary philosophy continuously confuses these two very different sorts of relations. Naturally the question arises of how it is possible to surmount our relation to the object so as to determine whether objects themselves possess the properties we encounter in relating to objects. In other words, given that we can only ever relate to the object in relating to the object how is it possible to surmount this relation to get at the being of the object itself? Much more will have to be said about this later– and the answers will be surprising with respect to standard prejudices about realism –however, for the moment it can be said that onticology takes its epistemological inspiration from the transcendental realism of Roy Bhaskar. Among other things, Bhaskar sought to provide a transcendental grounding for the sciences. Insofar as onticology defends the thesis that the field of being is much more vast than the field of objects investigated by the natural sciences, it parts way with the thesis that the domain of being is exhausted by the domain of natural objects. However, the general form of Bhaskar's argument holds for our realist purposes. A transcendental argument seeks to elucidate the conditions under which certain acknowledged practices and forms of cognition are possible. Kant, for example, asked what must be the case for mathematical judgments to be possible. How is it both that we are able to extend our knowledge, as if by magic, through mathematical judgments and, more significantly, that these judgments are able to provide genuine knowledge of the world despite the fact that these

forms of reasoning are not based on experience? Part of Kant's argument consisted in claiming that mind imposes the forms of space and time on the data of experience. In other words, space and time are not attributes of being itself but rather of the mind that regards being. Insofar, Kant argues, as mathematics is ultimately a rumination on the nature of space and time taken in their most abstract form and insofar as the mind imposes space and time on the manifold of sensation, it thus follows that a priori judgments about the nature of spatio-temporal relations are possible that anticipate the structure of actual-space times without directly experiencing these space-times. Why? Because any manifold of sensation must necessarily be structured by these forms imposed by intuition. **Onticology– A Manifesto for Object-Oriented Ontology Part I Posted by larval subjects under Object-Oriented Philosophy My Question Is This: You Say, "Granted There Are Deeper And Deeper And More Primitive Forms Of "Knowledge" Or Abstraction Into Mechanism, But I Wouldn't Say That The Atomic Level Brings A New Epistemic World Into View." But Doesn't Any Kind Of Knowledge, However Primitive, Imply Bringing Some Kind Of Epistemic World Into View Or Proto-View?** Joel: David, Yes Exactly. My Question Is, Just How Far Down The Complexity Gradient Can The Meaning Of 'Knowledge,' Or 'Semiotics, Or Specifically Representation (Maya) Be Stretched Without Breaking? I Can Comfortably Stretch It To The Genetic Level, As The Proto-Epistemic, Where Genetic Codes 'Represent' Phenotypic Structures Or Modes, And The Whole Of Evolution Learns From Its Mistakes And Successes As That Knowledge And History Is Encoded Into And As The Recapitulation Of The Embryogenesis Of The Organism. But An Atom, In My Model, While Certainly Reacting To Its Environment From Its Own Attractor, Essence, And 'Interiority' (Prehension) Isn't Making Representations Of It To Do So. But The Key To The Power And 'Essence' Of Representation (And Simultaneously Its Limitation) Is The Capacity Of Abstraction And Choice. Evolution Does This At A Rudimentary Level (And Anthropomorphizing A Bit) With Its Sending Into The World Its 'Random' Variations, Organismic Ideas, In A Sense. Its 'Mind' Is Literally On The Outside As Our Living World And Biosphere. A 'Choice' Is Made In The Biospheric Interiority,

And Knowledge For The Evolutionary/Embryogenetic Trajectory Is Gained By The "Differential Reproductive Success" Of These Individuals Lives Through Time, As The Increasing Intelligence, Representation, And Capacity For Choice Is Directly Injected Into The Flow Through Sexual Selection. But Where Is The Encoded/Abstracted (Enfolded) Representation Of Options And A Choice Among Them Made At The Atomic Level? I Am Working With The Most Cutting-Edge And Coherent Models For The Atom (Including My Own) That I Can Find Among The Heterodoxy, And I See No Evidence Of Representation And Choice At This Level, Or Rather No Real Way To Further Stretch The Meaning Of The Terms To This Level. Sure There Is Uncertainty From Infinite Difference And Immanent Causation (See Bohm On Infinite Causation), But Without Representation There Really Is No Choice (Although Crudely We Could Call A 'Bifurcation' A 'Choice' By The System As A Whole). So I Am Simply Saying That Representation At The Atomic Level Is Entirely Enfolded And These Are Really Just Ontic Level Phenomena. The Atomic Level, In My View, Is Pre- Not Yet Proto-Epistemic, Which I Place at the Genetic Level Where We Find the First Codes And (e&eb) Primitive Representation. This Just Means That There Is A Real Distinction Between Ontic And Epistemic (And A Gradient Between Them) And We Can't Collapse The Two Into Just The Ontic-Epistemic. And Yet, With The Distinction Intact, And The Gradient To Explore, They Remain Nondual, Univocal, ONE Or Ontic. What I Would Say Exists At All Levels (My Own Self-Aware Myth About The Given) Is What I Call The "Symbiogenesis Of Subject And Object". This Is The Spinozan Essence Of Dynamic Stability And Growth, And Leibnizian Prehension. Also Called The "Nucleation Of Observability" In Spinbitz. It Is From This Nucleus Of The Observer (Fuller) In A "Point-Free Geometry" (Whitehead) Of Pointless "Points Of View" As Conditions Of Boundary That Wilber's "All Is Perspective" Finds Its Ontic Grounding. And This Symbiogenesis Is Also A Rudimentary Or Deep-Level Native Or Primitive Intelligence, Evolution At The Involution Of The Subject-Object Interface. It Is An Exploration Of Creative Learning And Direct Immanent Awareness, Unmediated (Enfolded) By Representational Forms (Unfolded). It Just Is, In Waves (Cosmic Vertebrae)

Before It Abstracts, Represents And Incarnates The Flesh Of What Could Be. Matter Is Enfolded Seed Of The Flower Of Its Abstraction. Maya Is This Depth And Recursion Of Brahma Into And Through The Boundary Conditions Of Abstract Relation. And Brahma Is The Span Of The Emergent Rigor Of The Cosmic Ergodic Spine That Opens The Recursions Into The Field Of Infinite Difference. In Immanent-Transcendent Waves, Literally And Empirically A Limit-Cycle In Ergodicity, It Enfolds And Unfolds From Anatom To Anatomy, From Infinite Intelligence Or Omnirelational Awareness Of Self, Through The Representational Depths Of Self As Self-Consciousness. Omni-Evolution, Trans-Dynamics, And The Anatom Posted On February 21, 2013A Discussion With Tom Huston, Part II The Following Is Part 2 Of An Excerpt From A Discussion Initiated By Tom Huston On His Facebook Page, Edited And Embellished For Clarity. Find Part One, Here. In This Part We Dig Into The Nature Of The "Anatom" And The "Ergodic Spine" Of An Electro-Fractal Cosmos, And Into The Trans-Dynamic Integration Between Being And Becoming, And The X-Interface−The Crossroads Of The Ontic-Epistemic (Brahma And Maya) And The Subject-Object Polarities. Univocal_Dynamics_V1 Tom Huston Is A Founding Member Of Integral Institute And A Former Editor Of Enlighten Next Magazine. His Webpage Is Tomhuston.Com. David Marshall Is A Writer And Editor Living In Chicago. No meaning, value, or normativity to be found in nature, that there is nothing **natural beings ought to be,** but that, rather, these judgments arise from us (See last post of larval subjects) which has led some to raise valuable questions about the coherence of these claims. The problem arises when the following three propositions are taken together: 1. there is nothing outside of nature.2. Beings have no intrinsic meaning, purpose, or value (in and of themselves, there's nothing they ought to be). 3. Value judgments about what beings and being ought to be arise from us and beings like us (bonobo apes, dolphins, institutions, birds of paradise, etc). The problem arises between proposition 1 and 3. How can it both be true that there is nothing outside of being and those normative judgments belong to us and other beings capable of making normative judgments, not nature? The problem arises from restricting these judgments to humans and beings capable of

making these judgments. In making such a claim it seems as if we're saying that there's something outside of nature, something that is beyond nature, thereby violating the first thesis and potentially reintroducing the nature/culture distinction. By contrast, if we say that normative judgments are the special domain of those living beings with the proper degree of sentience to make such judgments, then we seem to reintroduce the nature/culture distinction and fall back into the sorts of problems that I outlined in my last post (and that are so nicely critiqued by thinkers such as Latour). Is there a way out of this? I don't know. Great_Chain_of_Being_2_ (lighter) ASIDE: It seems to me that this is really what the debate between realism and anti-realism, realism and socio-linguistic constructivism surrounding the new materialisms and speculative realists has really been all about.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \qquad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9}+\frac{N_9 dG_9}{dR_9}-\frac{dB_9}{dR_9}>\frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9}+\frac{dM_9}{dR_9}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usance bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_9=\frac{N_9 G_9}{2M_9}\frac{d(9nF_9)}{dR_9}(\geq 0) \qquad (9a)$$

$$I_9=-\frac{N_9 G_9}{2M_9}\frac{d(9nM_9}{dR_9}(\geq 0) \qquad (9b)$$

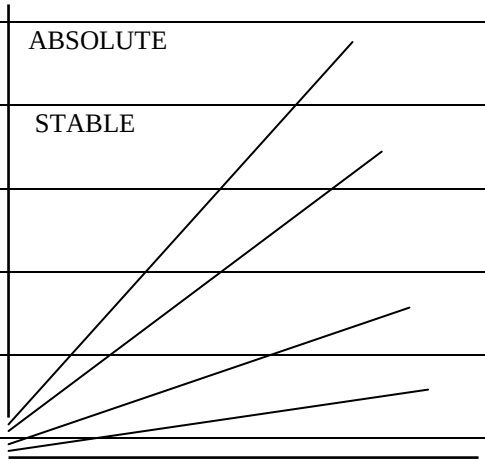
It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9>L_9-I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_9>L_9-I_9$, when $P_9^2/4>I$ (8a)

$P_9>L_9-I_9+\frac{(I_9-P_9^2}{(L_9-I)}$ when $P_9^2/4<I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Reffer couffesy Nunney

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability

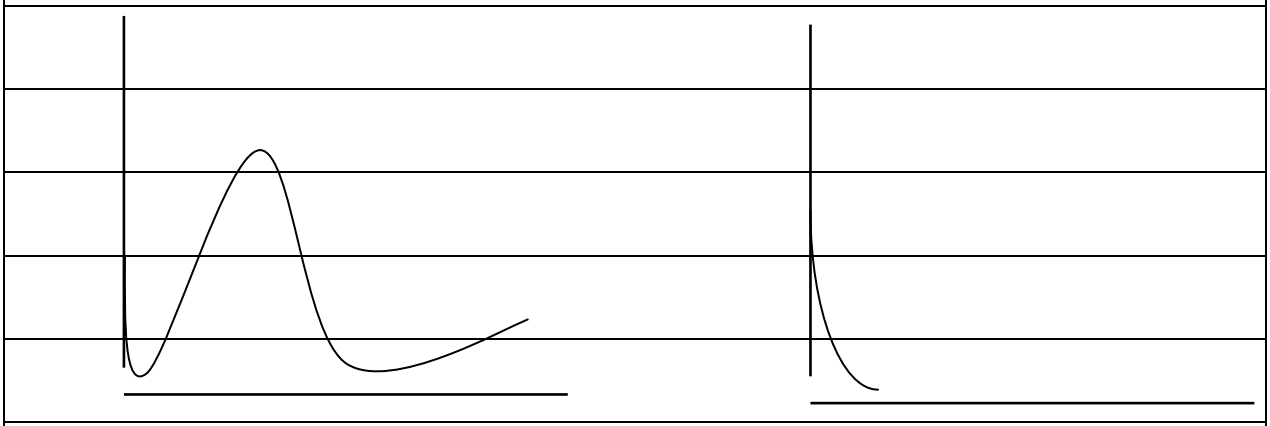
 

is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

Indeed, Nunny's model can be applied to some of the consciousness and philosophical problems. Predator prey analysis has ripped open the veneer of the fundamental concept of universe namely Debit-Credit world.

The noise of ordinary consciousness and (e&eb) beliefs drowns out (e) the signal. Unconsciousness is the background of (e) our ordinary awareness. Our organism is very much at (eb) the center of such effects. The organismic source is (=) our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is a meme-Scape. Pre-conceived concepts vie with structures, concepts with images. Like scientists who ignore assumed truths, we leapfrog over our beliefs and personality deficits, claiming idiosyncratic imagination is literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is both. We inherit the structures that make our experience what it is.

But the structure itself is “empty,” and each human culture “fills” it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us

feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through 're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns, nomenclature, and internal coherence of meaning and representation. In a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist

simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** The lower domain of experience is what we do not directly observe, the quantum realm, while the higher level is what we ordinarily observe, the classical realm. Based purely on experiment, the formulation of quantum theory initially placed the great divide between the quantum and classical domains between the measuring apparatus and the particle (Herbert, 1985). However, in the early 1930s, John von Neumann, in his rigorous mathematical treatise on quantum mechanics, found no support for such a division, compelling him to conclude that the wave for was collapsed by consciousness. Von Neumann's treatise has been called the "quantum bible" and "the most influential book on quantum theory ever written" (Herbert, 1985).

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t-T_9)F_9\big(R_9(t-T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t-T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ and the effective prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting $dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

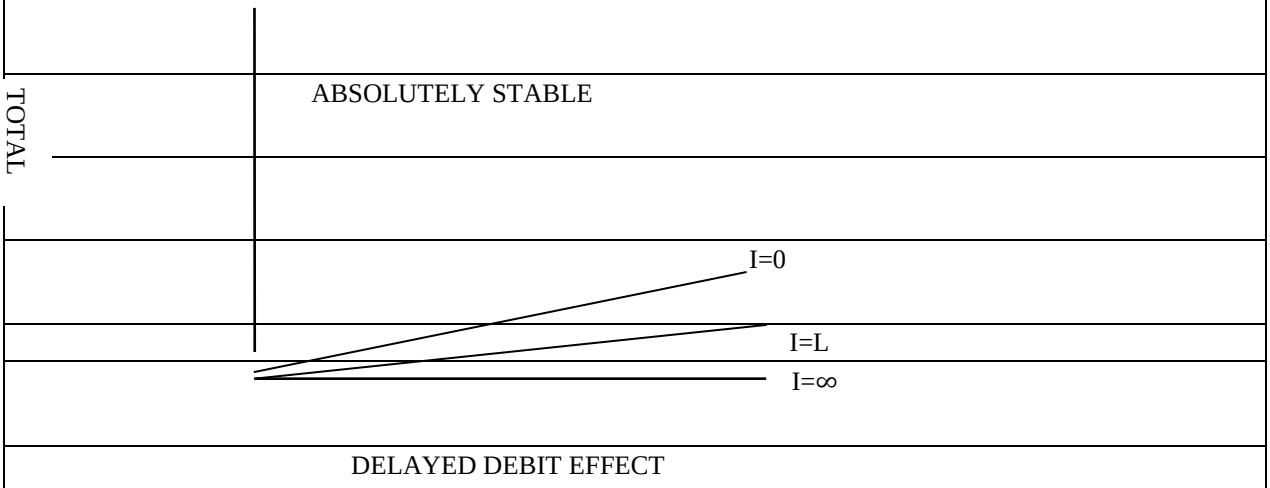



FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.

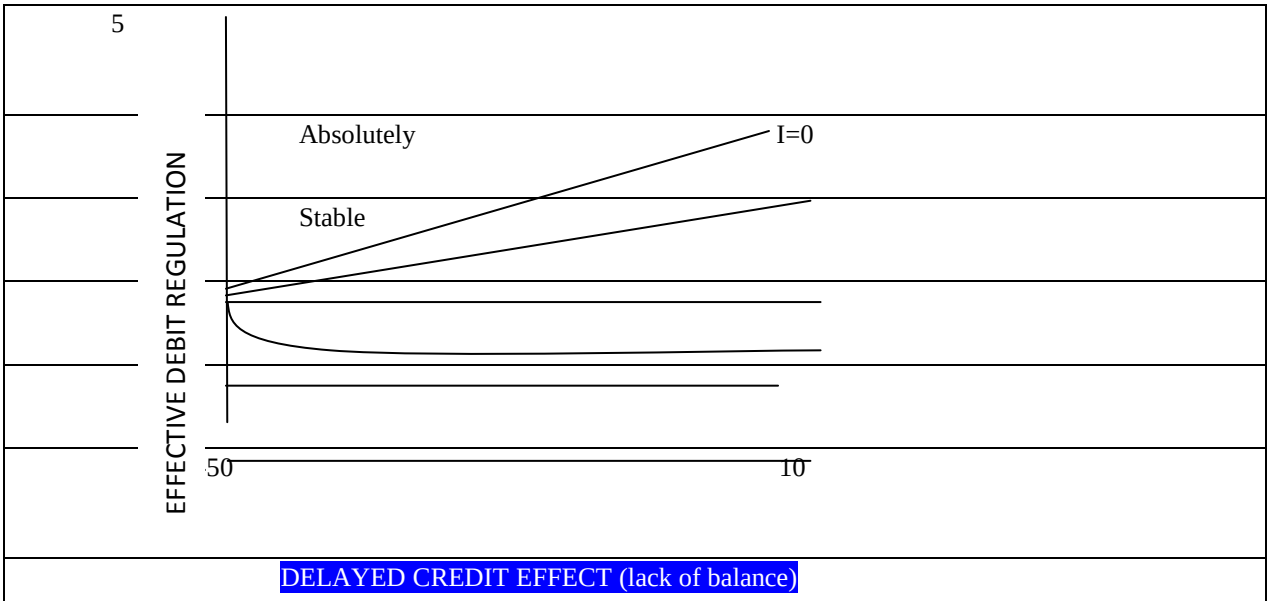



Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the

constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^\infty \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9\left(U_9 - V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9\right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$ then given the constraints of $f(T_9)$, $(A1.5)$ *can never be satisfied* Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDERATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 - V_9M_9 -]e^{-T_9\lambda_9} = 0$$

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9 Z_9 \sin(Z_9 T_9) - (A_9 M_9 + Y_9) + (A_9 M_9 + X_9)\cos(Z_9 T_9) = 0$$

$$(A_9 M_9) Z_9 - M_9 Z_9 \cos(Z_9 T_9) + (A_9 M_9 - X_9)\sin(Z_9 T_9) = 0$$

Where $A_9 = U_9 - V_9$ Using (A2.3), substitute for $Z_9^2$ and $Z_9$ in (A2.2):

$W_9^2(X_9^2 - 3M_9 A_9 X_9 - M_9^2(X_9 - Y_9) + 2M_9^3 A_9 + 2A_9^2 M_9^2]$

$$+W\left[MA^{3}-A^{2}\mathrm{X}+2\mathrm{MA}(2\mathrm{X}+\mathrm{Y})-2X^{2}\right]+A^{2}(\mathrm{X}+\mathrm{Y})=0$$

*where* $W_9$=1-cos ($\mathrm{Z}_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (\mathrm{X}_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ \mathrm{X}_9 - Y_9/2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range 0⍰$W_9$ ⍰2 *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (\mathrm{X}_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\,\mathrm{M}_9)^2\,(2A_9\,\mathrm{M}_9 - (\mathrm{X}_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9M_9^2 - 3M_9X_9\right) + W_9\left(3M_9A_9^2 - 2A_9X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9 + Y_9\right)$$

$=2M_9^2W_9^2(M_9 + A_9) + M_9A_9^2W_9 + \left(2A_9M_9 - X_9 + Y_9\right)(A_9W_9 + M_9W_9^2$

$+ (2- W_9)((2M_9X_9 + M_9Y_9)W_9 + A_9(X_9 + Y_9))$,

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range 0≤ $W_9$ ≤2.Defining

$Q_9(W_9, A_9) = a_9W_9^2 + c_9$ *where* $a_9$ $b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

**Harald Atmanspacher, Hans Primas** show that such an epistemic perspective can be consistently complemented by Einstein's ontically oriented position. **Time, Quantum and Information 2003, pp 301-321 Epistemic and Ontic Quantum Realities Harald Atmanspacher, Hans Primas Robert W. Spekkens** present a toy theory that is based on a simple principle: the number of questions about the physical state of a system that are answered must always be equal to the number that are unanswered in a state of maximal knowledge. Many quantum phenomena are found to have analogues within this toy theory. These include the noncommutativity of measurements, interference, the multiplicity of convex decompositions of a mixed state, the impossibility of discriminating nonorthogonal states, the impossibility of a universal state

inverter, the distinction between bipartite and tripartite entanglement, the monogamy of pure entanglement, no cloning, no broadcasting, remote steering, teleportation, entanglement swapping, dense coding, mutually unbiased bases, and many others. The diversity and quality of these analogies is taken as evidence for the view that quantum states are states of **incomplete knowledge** rather than states of reality. A consideration of the phenomena that the toy theory **fails to reproduce, notably, violations of Bell inequalities** and the existence of a Kochen-Specker theorem, provides clues for how to proceed with this research program.DOI: http://dx.doi.org/10.1103/PhysRevA.75.032110 **Evidence for the epistemic view of quantum states A toy theory Phys Rev A 75, 032110 – Published 19 March 2007 Robert W. Spekkens In such specialized and autonomous areas, it is to the eminent and erudite authors to interpret the contradistinctions and contradictions' in the dove tailing explaining. This I feel would be yeoman service in the creation of Nature's general ledger. M. S. Leifer and O. J. E. Maroney** examine the relationship between quantum contextuality (in both the standard Kochen-Specker sense and in the generalized sense proposed by Spekkens) and models of quantum theory in which the quantum state is maximally epistemic. We find that preparation noncontextual models must be maximally epistemic, and these in turn must be Kochen-Specker noncontextual. This implies that the Kochen-Specker theorem is sufficient to establish both the impossibility of maximally epistemic models and the impossibility of preparation noncontextual models. The implication from preparation noncontextual to maximally epistemic then also yields a proof of Bell's theorem from an Einstein-Podolsky-Rosen-like argument. DOI: http://dx.doi.org/10.1103/PhysRevLett.110.120401 Received 25 August 2012Published 20 March 2013 © 2013 American Physical Society **Maximally Epistemic Interpretations of the Quantum State and Contextuality Phys. Rev. Lett 110, 120401 – Published 20 March 2013 M. S. Leifer and O. J. E. Maroney.** Epistemological, onitic and ontological aspects particularly of quantum mechanics have invited

profound thinkers to contribute their mite to the growing literature on the subject, which includes the contributions from the subterranean realm and ceratoid dualism of philosophy like that Chaos theory of Deleuze, Derrida , Kierkegaard and others. **Andrei Khrennikov** shows that the Dirac-von Neumann formalism for quantum mechanics can be obtained as an approximation of classical statistical field theory. This approximation is based on the Taylor expansion (up to terms of the second order) of classical physical variables – maps f: Ω → R, where Ω is the infinite-dimensional Hilbert space. The space of classical statistical states consists of Gaussian measures ρ on Ω having zero mean value and dispersion σ2 (ρ) ≈ h. This viewpoint to the conventional quantum formalism gives the possibility to create generalized quantum formalisms based on expansions of classical physical variables in the Taylor series up to terms of nth order and considering statistical states ρ having dispersion σ2(ρ) = h n (for n = 2 we obtain the conventional quantum formalism). **December 2005, Volume 18, Issue 7, pp 637-650 Date: 28 Nov 2005 Generalizations of Quantum Mechanics Induced by Classical Statistical Field Theory Andrei Khrennikov** It is widely accepted that consciousness or, in other words, mental activity is in some way correlated to the behavior of the brain or, in other words, material brain activity. Since quantum theory is the most fundamental theory of matter that is currently available, it is a legitimate question to ask whether quantum theory can help us to understand consciousness. Several approaches answering this question affirmatively, proposed in recent decades, will be surveyed. It will be pointed out that they make different epistemological assumptions, refer to different neurophysiological levels of description, and adopt quantum theory in different ways. For each of the approaches discussed, these imply both problematic and promising features which will be indicated. **Discrete Dynamics in Nature and Society Volume 2004 (2004), Issue 1, Pages 51-73 http://dx.doi.org/10.1155/S102602260440106X Quantum theory and consciousness: an overview with selected examples Harald Atmanspacher Copyright © 2004 Hindawi Publishing**

**Corporation. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.** Another thought provoking and investigatory article is by Peter G. Lewis et al. Perhaps the quantum state represents information about reality, and not reality directly. Wave function collapse is then possibly no more mysterious than a Bayesian update of a probability distribution given new data. **Peter G. Lewis et al** consider models for quantum systems with measurement outcomes determined by an underlying physical state of the system but where several quantum states are consistent with a single underlying state—i.e., probability distributions for distinct quantum states overlap. Significantly, they demonstrate by example that additional assumptions are always necessary to rule out such a model. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.150404 **Distinct Quantum States Can Be Compatible with a Single State of Reality Phys. Rev. Lett. 109, 150404 – Published 9 October 2012 Peter G. Lewis et al** Communication complexity of a quantum channel is the minimal amount of classical communication required for classically simulating a process of state preparation, transmission through the channel and subsequent measurement. It establishes and perpetuates a limit on the power of quantum communication in terms of classical resources. **Alberto Montina** shows that classical simulations employing a finite amount of communication can be derived from a special class of hidden variable theories where quantum states represent statistical knowledge about the classical state and not an element of reality. This special class has attracted strong interest very recently. The communication cost of each derived simulation is given by the mutual information between the quantum state and the classical state of the parent hidden variable theory. Finally, find that the communication complexity for single qubits is smaller than 1.28 bits. The previous known upper bound was 1.85 bits. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.110501 **Epistemic View of Quantum States and**

**Communication Complexity of Quantum Channels Phys. Rev. Lett 109, 110501 – Published 12 September 2012 Alberto Montina for models see elsewhere .** Stability analysis of what? Quantum state representing reality or our knowledge of reality? Does the quantum state represent reality or our knowledge of reality? In making this distinction precise, **Nicholas Harrigan, Robert W. Spekkens** are led to a novel classification of hidden variable models of quantum theory. They show that representatives of each class can be found among existing constructions for two-dimensional Hilbert spaces. Approach also provides a fruitful new perspective on arguments for the nonlocality and incompleteness of quantum theory. Specifically, **Nicholas Harrigan, Robert W. Spekkens** show that for models wherein the quantum state has the status of something real, the failure of locality can be established through an argument considerably more straightforward than Bell's theorem. The historical significance of this result becomes evident when one recognizes that the same reasoning is present in Einstein's preferred argument for incompleteness, which dates back to 1935. This fact suggests that Einstein was seeking not just any completion of quantum theory, but one wherein quantum states are solely representative of our knowledge. Hypothesis is supported by an analysis of Einstein's attempts to clarify his views on quantum theory and the circumstance of his otherwise puzzling abandonment of an even simpler argument for incompleteness from 1927. **February 2010, Volume 40, Issue 2, pp 125-157 Date: 09 Jan 2010 Einstein, Incompleteness, and the Epistemic View of Quantum States Nicholas Harrigan, Robert W. Spekkens** statistical model of the probabilistic description of physical reality has been studied by many authors. **Andrei Khrennikov posits** a contextual statistical model of the probabilistic description of physical reality. Here contexts (complexes of physical conditions) are considered as basic elements of reality. There is discussed the relation with QM. He proposes a realistic analogue of Bohr's principle of complementarity. In the opposite of the Bohr's principle, **Andrei Khrennikov** principle has no direct relation with mutual exclusivity for

observables. To distinguish his principle from the Bohr's principle and to give better characterization, he changes the terminology and speaks about supplementarity, instead of complementarity. Supplementarity is based on the interference of probabilities. It has quantitative expression trough a coefficient which can be easily calculated from experimental statistical data. There is need not appeal to the Hilbert space formalism and noncommutativity of operators representing observables. Moreover, in our model there exists a pair of supplementary observables which cannot be represented in the complex Hilbert space. There are discussed applications of the principle of supplementarity outside quantum physics. **October 2005, Volume 35, Issue 10, pp 1655-1693 The Principle of Supplementarity: A Contextual Probabilistic Viewpoint to Complementarity, the Interference of Probabilities and Incompatibility of Variables in Quantum Mechanics Andrei Khrennikov Foundations of Physics** Description of the vacuum in Yang-Mills theory and arrive at a physical interpretation of the pseudoparticle solution and the attendant violation of symmetries has been done by **R. Jackiw and C. Rebbi.** It is here vaccuum stability forms the bastion, pillar, post and stylobate of the mechanisms of yang mills Theory and action of a particle. The existence of topologically inequivalent classical gauge fields gives rise to a family of quantum mechanical vacua, parametrized by a CP-nonconserving angle. The requirement of **vacuum stability** against gauge transformations renders the vacua chirally noninvariant. DOI: http://dx.doi.org/10.1103/PhysRevLett.37.172 **Vacuum Periodicity in a Yang-Mills Quantum Theory Phys. Rev. Lett 37, 172 – Published 19 July 1976 R. Jackiw and C. Rebbi.** Three-dimensional Yang-Mills and gravity theories augmented by gauge-invariant mass terms are analyzed. These topologically nontrivial additions profoundly alter the particle content of the models and lead to (eb) quantization of a dimensionless mass-coupling-constant ratio. The vector field excitations become (=) massive, with (e&eb) spin 1 (rather than massless with spin 0), and the mass provides (eb) an infrared cutoff. The gravitation

acquires (e) mass, mediates (e&eb) finite-range interactions, and has (e) spin 2 (rather than being absent altogether); although its mass term is of (e)third derivative order, there are no ghosts or acausalities. DOI: http://dx.doi.org/10.1103/PhysRevLett.48.975 © 1982 The American Physical Society **Three-Dimensional Massive Gauge Theories Phys. Rev. Lett 48, 975 – Published 12 April 1982 S. Deser et al** All scattering amplitudes in (eb) the maximally supersymmetric N=4 super-Yang–Mills theory possess (e) a new, dual superconformal symmetry which extends the previously discovered dual conformal symmetry of MHV amplitudes. To reveal this property authors formulate the scattering amplitudes as (=) functions on (e&eb) the appropriate dual superspace. Rewritten in this form, all tree-level MHV and next-to-MHV amplitudes exhibit manifest dual superconformal symmetry. Authors propose a new, compact and Lorentz covariant formula for (e) the tree-level NMHV amplitudes for (e) arbitrary numbers and types of external particles. The dual superconformal symmetry is broken at loop level by infrared divergences. However, authors provide evidence that the dual conformal anomaly of (e) the MHV and NMHV superamplitudes is (=) the same and, therefore, their ratio is (=) dual conformally invariant. Authors show this explicitly for (e) the six-particle amplitudes at one loop. Authors conjecture that these properties hold for all, MHV and non-MHV, superamplitudes in N=4 SYM both at (eb) weak and at strong coupling. **Nuclear Physics B Volume 828, Issues 1–2, 21 March 2010, Pages 317–374 Dual superconformal symmetry of scattering amplitudes in N=4 super-Yang–Mills theory J.M. Drummonda DOI: 10.1016/j.nuclphysb.2009.11.022.** By comparison with numerical results in (eb) the maximal Abelian projection of lattice Yang–Mills theory, it is argued that the nonperturbative dynamics of (e) Yang Mills theory can be described by (e) a set of fields that take their values in (eb) the coset space SU(2)/U(1). The Yang–Mills connection is parameterized in (eb) a special way to separate (e&eb) the dependence on the coset field. The coset field is then regarded as (=) a collective variable and a method to obtain its effective

action is developed. It is argued that the physical excitations of (e) the effective action may be (=) knot solitons. A procedure to calculate the mass scale of knot solitons is discussed for (e) lattice gauge theories in the maximal Abelian projection. The approach is extended to (e&eb) the SU (N) Yang–Mills theory. A relation between the large N limit and the monopole dominance is pointed out. **Physics Letters B Volume 458, Issues 2–3, 8 July 1999, Pages 322–330 An effective action for monopoles and knot solitons in Yang–Mills theory Sergei V. Shabanov DOI: 10.1016/S0370-2693(99)00612-7.** It is here the quintessential formulation of consciousness occurs. Like a man standing on the threshold of infinity trying to ponder what lies beyond the veil, which separates the seen from unseen a true seeker tries to lift the individual general ledger to the cosmic general ledger. Like in Darwinian evolution, mind also has to evolve from individual consciousness to cosmic consciousness. In Kaishmiri Shaivism one withdraws from the finite to the infinite, but one also goes on an outward journey from the infinite to the finite, because both the finite and the infinite have an intimate connection. The finite is not seen as unreal, but as a symbol of the infinite. There is no real distinction between them. Those two movements constitute Spanda, a key concept in Kaishmiri Shaivism. Spanda is the pulsation of the Absolute in different phases of being. There are no opposites like subject and object, unity or duality, absolute or relative. They are just different phases of the universal vibration of the Absolute. The goal is to realize or be (=) at once infinite and finite. One does not turn away from appearances (like in Advaita Vedanta), but one realizes that the Absolute manifests all things. Spanda, the eternal pulsation of the Absolute, oscillates between a passion to create and dispassion from the created. Through it the Absolute transform itself into all things and then returns back into the emptiness of its undifferentiated nature. In Kaishmiri Shaivism the Absolute is seen as pure consciousness (=being). The Absolute is an eternal all-pervasive principle, the highest reality, the nature of all entities eternally and blissfully at rest within its own nature. The Absolute is the nature of the Self (and

thus of us all). The Absolute is divine; it is Shiva, the Lord of the Universe. It is full of (e) conscious activity through (e&eb)which it generates the universe, and reabsorbs it into (e&eb) itself at the end of each cycle of creation. Thus we speak of monism, as everything resides within (eb) this one absolute consciousness. It sustains (eb) all things, it embraces (e) all things, and it pervades (e&eb) all things. All things are appearances within (eb) the absolute consciousness, but nevertheless real (in contrast to (e) Advaita Vedanta where (e) appearances are seen as unreal or illusionary). All things appear external (out there, outside ourselves), but they do not have (e) a being on their own. They do not exist as (=) separate entities on their own. Everything is contained within (eb) consciousness. What we see as objects are (=) manifestations of consciousness. The event which constitutes (eb) the universe is always internal events happening within (e&eb) consciousness because (e) their essential nature is consciousness itself. If a physical object were totally material, and independent or external to consciousness, it could never be experienced. The universe and (e&eb) consciousness are two aspects of a whole. The universe is (=) an attribute of consciousness which bears consciousness as (=) its substance. Doctrine of Vibration: S.G.Dyczkowski what you "see" is (=) not always what you "get." Many people mistakenly take their own visions literally without (e) "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in (eb) the physical sense but in a psychological sense. The archetypes of (e) the collective unconscious are arrayed behind (e) the scenes of current worldwide conditions, of (e) crisis and confusion. They mirror (e&eb) our own states back at us, whether we perceive them as (=) such or interpret them plausibly or not (Note the essay has been written in model able form and models are given in various papers.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4\,a_9 \square - b_9 \square 0$$

And

$$b_9^{\,2} \square 4\,a_9 c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9\, in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\, X_9$ ,

$Let\ Y_9 = \ y_9\, X_9$And try the substitution $A_9 = X_9/(2 + A_9\ y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + \ n_9 y_9)$ we have

$$W_9^{\,2}\, y_9 \left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^{\,2}\left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9\, X_9 \left[\frac{X_9 (1 + n_9 y_9)}{M_9^{\,2}(2 + n_9 y_9)^2} + 2\, y_9 \left(n_9 - 1\right) + \frac{X_9^{\,2}(9 + y_9)}{M_9^{\,2}(2 + n_9 y_9)} = 0.\ (A2.12)\right.$$

Since $n_9$ and $A_9$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9 = (2 + n_9 y_9)$, *the condition*(A2.99?)*can be written as*

$[\frac{X_9(E_9-9)}{E_9^2}]^2[\frac{X_9^2}{M_9^2} - 4\, y_9\ E_9^2 \frac{X_9}{M_9^2} + 4y_9 (n_9 - 2)E_9^3]$ .

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9 / W_9^2)$ can be expressed as

$2\, E_9 [y_9 E_9 + (y_9\ E_9 (9 + y_9)(E_9 - n_9))^{9/2}] > \frac{X_9}{M_9^2} > 2E_9 [y_9 E_9 - y_9 E_9 (9 + y_9)(E_9 - n_9))^{9/2}]$

Concept Of Being From Shaivite Point Of View: In His Famous Work Entitled On The Conservation Of Force (1847), Hermann Von Helmholtz Had Divided The Totality Of The Energy In The Universe Between Potential Energy And Kinetic Energy And Affirmed The Reciprocal Convertibility Of The Two. In 1852, William Thomson Pointed Out That There Exists A Sub-Ensemble Within Kinetic Energy, Heat, Which, Once It Has Been Generated, Is No Longer Entirely Convertible Into Potential Energy-Or Into Any Other Form Of Kinetic Energy. Considering That The (Partial) Reconversion Of Heat Into Labor Is Possible Only In Situations That Present A Disparity In Temperature, And That Heat Tends To Pass From Warmer To Cooler Bodies By Spreading On An Even Temperature Level Through Space, Thomson Concluded That The Universe Tends Towards A Final State Where Any Energetic Transformations, Every Movement And Every Form Of Life Will Cease: The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO: The Ultimate Nature Of Anything Cannot Change. Change Of Any Sort Is Merely Apparent (Vivarta). The World Of Change And Becoming Is A False Super- Imposition (Adhyaropa, Adhyasa) On The Absolute. It Appears So On The Screen Of Individual, Collective Consciousness, Albeit One Cannot Be So Sure Of Cosmic, For There Are Many Reasons That Cosmic Transactions Also Had To Be Subjected To Change

By ParaBrahaman. In Cosmic Terms, The Mistake (Bhranti) Consists Of The Supposition That The Real Brahman Is The Unreal Universe And The Unreal Universe Is The Real Brahman. In Microcosmic Terms, It Is The Mistake Of Falsely Conceiving The Body, Mind Or Even One's Personality To Be The Self. Note We Have Taken Self As Witness Consciousness. Notwithstanding There Is Nothing Wrong In The Conception Of Self As A Dynamical Presupposition, In That Individual Consciousness And Collective Consciousness And Even Cosmic Consciousness Evolutes In To ParaBrahaman. In The Same Way As The Image Of A Snake Is Falsely Superimposed On A Rope, Similarly The Universe Is Falsely Projected Onto The Real Substratum, The Brahman. Ignorance Is Not Merely A Personal Lack Of Knowledge, But A Cosmic Principle. As Such It Is Called "Maya," The Indefinable Factor (Anirvacaniya) That Brings This Mistake In Identity About. The Reality Status Of This Cosmic Illusion Is Also Indefinable: On The One Hand It Is Not Brahman, The Sole Reality; On The Other Hand It Is Not Absolutely Non-Existent Like A Hare's Horn Or The Son Of A Barren Woman. Note Here The Vedantin Avoids Talking About Zero Which We Have Avidly And Assiduously Stressed About. Nothing Is The Factor That Exists. Also See the Remarks about Satre's Being And Nothingness. Brahman Is The Source Of World Appearances Only In The Sense Of Being Their Unconditioned Ground Or Essential Nature. The Universe Is False Not Because It Has No Nature Of Its Own But Because It Does Have One. The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski. Concept Of Being From Shaivite Point Of View: It Is The 'I Am The Body/Mind' Belief That Gives Rise To The 'I Am Not The World' Belief. These Two Beliefs Are (=) Co-Created."Consciousness Projects The Appearance Of The Mind, Body And World By Taking (Eb) The Shape Of Thinking, Sensing And Perceiving."“Attention Is Consciousness With An Object. When The Object Vanishes, Attention Simply Remains What It Always Is Consciousness."“There Is No Purpose To Meditation. The Purpose Is Already Accomplished."“Everything

That Is Experienced Is Experienced By, Through, In And As Consciousness.”“The Seen Cannot Be Separated From Seeing And Seeing Cannot Be Separated From Consciousness.” “**The Reality Of Any Experience Is Not Hidden In The Appearance; It Is Expressed By The Appearance The Transparency Of Things**' By Rupert Spira Book Review By Dennis Waite The Following Is A Review Of Rupert's Book: 'The Transparency Of Things: Contemplating The Nature Of Experience’. There Is An Essay From The Book Here And The Essay On Cézanne May Be Read At The Awakened Eye Website Nature Rejects The Naiveté That Seeks Absolute Truth. We Are Beginning To Realize, Individually And Culturally, That "Realities" Are All Human Constructions. The Task Becomes One Of "Catching Ourselves In The Act" Of Creating Our Own "Reality" From The Flow Of Events. Human Truth Is Always An Engagement Of Mind With Experience. The Challenge Of The Therapist In These Times Of Chaotic Change Is To Validate The Concept That We Don't Have To Fear The Collapse Of What We Think We Are. Strange Attractors: Transference, Holography, And An **Archetype Burke, J. (2003). Strange Attractors: Transference, Holography, And An Archetype (Doctoral Dissertation, Pacifica Graduate Institute, 2003).** When we get to the topic and understanding of the Transcendental Logic, the brilliant and baffling Transcendental Deduction (TD) Recall the two movements just discussed, the one from experience to its conditions and the one from the forms of valid inference to the concepts that we must use in all judging (the Categories). This duality led Kant to his famous question of right (quid juris) (A84=B116): with what right do we apply the Categories, which are not acquired from experience, to the contents of experience? (A85=B117). Kant's problem here is not as arcane as it might seem. It reflects an important question: How is it that the world as we experience it conforms to our logic? In briefest form, Kant thought that the trick to showing how it is possible for the Categories to apply to experience is to show that it is necessary that they apply (A97).TD has two sides, though Kant never treats them separately John Marlado,

accounts, attributes, and ascribes the philosophy of Nishida KitarÅ is lucid, insightful, and deeply informative despite the highly questionable argument that gives the book its structure. In comparison with other, technically more accurate accounts of Nishida's philosophy, it is broader in scope than most but more single-minded in its approach. The analysis is for the most part limited to four books that have been translated into English, out of approximately 16 volumes of philosophical essays in Japanese, but it aims to capture the contours and most important details of Nishida's entire philosophy. To appreciate the achievement of this account, one would need only to read, in English translation or in the Japanese original, any of the volumes with their streams of repetitive, puzzling expressions and circuitous, often dead- end argumentation. Wilkinson's elucidation of details in the passages he samples makes it easier to understand why Nishida is considered Japan's greatest academic philosopher. A critical look at the frame of the book and some of its own claims reveals where, to my mind, it has gone astray. The argument that frames this book is repeated throughout: Nishida's philosophy is an attempt to articulate, in a characteristically Western philosophical manner, the Zen experience he had as a young man. Nishida first made the attempt by using the conceptual frameworks of Western philosophers, and after they failed him he developed his own, largely by recasting Buddhist insights. In the end, a philosophical articulation of the Zen worldview in Western terms or even in Eastern terms recast by conceptual, Western philosophical methods approved impossible. What Nishida attempted is incommensurable with Western philosophy. While the author first offers these statements as hypotheses (p. 2), one premise of his argument that there is such a thing as "the Zen world-view" or "the Zen experience “remains unquestioned; and another that Nishida's pivotal experience was the Zen experienceâis asserted as "beyond question" (p. 151; see also pp. 28, 48, 10, 114).Wilkinson's presentation of Zen in the first chapter is drawn largely from D. T. Suzuki, with help from Suzuki's portage Abe Masao; some comments on time, birth, and death in; and a few other sources. The

result is a caricature that imagines Zen as an invariable, unitary experience. In fact, Dagan's teachings on practice as the manifestation of enlightenment do not sit well with D. T. Suzuki's emphasis on a satori experience, and Suzuki does not enlighten us about what Zen in its various practices has historically been. The author would have benefited from acquaintance with other presentations by Zen teachers, such as that of ShunryÅ Suzuki'sZen Mind Beginner's Mind, with historical accounts of Zen traditions, and with contrasts between DÅgen's view of time and Sam Â· sÄraand that of other Buddhist philosophers. A little research in contemporary Zen scholarship would have shown how implausible is the talk of "nirvana or awareness of mu" (p. 24), "the basic Zen position [that] never varies" (p. 3), "the Zen vision" (pp. 5, 7), "Zen epistemology" , "Zen experience" (p. 59), "the conception of time involved in Zen" (p. 15), and "the fundamental Zen assertion that there is no absolute time" (p. 16). Even if DÅgen's notion that time and being are nondual (p. 16) epitomized most of Zen, Nishida's early notion of "a transcendent, unchanging reality apart from time" (p. 51) would contrast sharply with it. Equally questionable is Wilkinson's contention that Nishida's "pure experience" is equivalent to Zen satori. The notion of pure experience Nishida developed in his first major work may be pivotal for his entire philosophy, but he never claims or implies that any part of his philosophy presupposes an exceptional experience like satori that grounds his convictions and renders them unverifiable to those who have not had the experience, as Wilkinson intimates (pp. 55, 159, and 161). Nishida presents his notions of pure experience and intellectual intuition, the grasp of its unity, and later... Perennialism and Constructivism is studied by Randolph T Dible : experiencer or the observer and the transcendental subjectivity itself Pure objectivity is essentially a deconstruction, apophatic and negativa The encountering of the experiencer or observer—transcendental subjectivity itself—at the foundation of the world leads inevitably to the recognition of pure objectivity as ultimate reality (which can be taken as its ultimate deconstruction, analogous to the apophatic or via

negativa), from which objects derive their value, weight, significance, meaning or objectivity. In this way, pure objectivity can be seen as the supra-self-evident Axiological Axiom, so to speak, even Unconditional Love, in romantic terms. This axiology (value theory) has a structure inverse to the relationship between transcendental subjectivity as the radical unity of pure self-reference and on the other hand, the world of forms, as mere traces (representations, indications) of the unique, original "first distinction" Spencer-Brown speaks of at the foundation of his calculus. That is, all forms (i.e., distinctions, differences) would reduce to being the first distinction, also known as the marked state, which can be called penultimate reality (pure self-reference or transcendental subjectivity: the Spirit which animates us), except that forms are complimentary to their content, which is their objectivity or value, which would reduce to the unmarked state or ultimate reality. It is the incongruity of form (thoughts; Whitehead's "negative prehensions") and value (feelings; Whitehead's "positive prehensions," or my notion of objectivity, meaning and qualia; in short, the non-formal aspects of experience) that holds forms open and keeps them from absolute reduction. This accounts for the brute, concrete persistence of the "functional illusion"-- to use a term from Dzogchen Buddhism-- of the world. Thus this system has an axiology of metaphysical objectivity grounded on the ideal of pure objectivity as the source of all value, meaning and significance, itself the very fecundity of profundity, which is the motive of drawing the distinction in the first place **(Wikipedia, Kant's writings and Logic Of Sense)** Nano rings were created by accident while intending to make quantum dots. They have interesting optical properties associated with excitons and the Aharonov–Bohm effect. Application of these rings used as light capacitors or a buffer includes photonic computing and communications technology. Analysis and measurement of geometric phases in mesoscopic rings is ongoing. Several experiments, including some reported in 2012, show Aharonov-Bohm oscillations in charge density wave (CDW) current versus magnetic flux, of dominant period h/2e through CDW rings up to 85 **μm** in

circumference above 77 K. This behavior is similar to that of the superconducting quantum interference devices (seeSQUID). **Aharonov–Bohm nano rings (Wikipedia) Stability analysis of nano rings would help progressively accentuate the research about the materials division that uses nano rings.** From the profound contemplation and wisdom of the Buddha and Mahabodhisattvas, we know there is no such reality. Instead, egolessness (non-self) is the only path to understand the reality of the deluded life. All existences are subject to the law of causes and conditions. These include the smallest particles, the relationship between the particles, the planets, and the relationship between them, up to and including the whole universe! From the smallest particles to the biggest matter, there exists no absolute independent identity. Egolessness (non-self) implies the void characteristics of all existence. Egolessness (non-self) signifies (eb) the non-existence of permanent identity for self and existence (dharma). Sunyata stresses (eb) the voidness characteristic of (e) self and existence (dharma). Sunyata and Egolessness possess (e) similar attributes. As we have discussed before, we can observe the profound significance of Sunyata from (e) the perspective of inter-dependent relationships. Considering dharma-nature and (e&eb) the condition of nirvana, all existences are immaterial and of (e) a void-nature. Then we see each existence as (=) independent of each other. But then we cannot find any material that does exist independent of (e) everything else. So Egolessness also implies (eb) void-nature! From the observation of all existences, we can infer (eb) the theory of nirvana and the complete cessation of (e) all phenomena. From the viewpoint of phenomena, all existences are so different from (e) each other, that they may contradict (e&eb) each other. They are (=) so chaotic. In reality, their existence is (=) illusionary and arises from (e) conditional causation. They seem to exist on one hand, and yet do not (e) exist on the other. They seem to be united, but yet they are so different to (e) one another. They seem to exist and yet they do cease! Ultimately everything will return to (e) harmony and complete calmness. This is the nature of all existence. It is the

final resting place for all. If we can understand this reality and remove our illusions, we can find this state of harmony and complete calmness. **Teachings in Chinese Buddhism (6) sunyata (emptiness) in the Mahayana context (Wikipedia) Many models have been given concatenating Buddhist doctrines and Vaccum energy. These models essentially throw light on the essence of self(witness consciousness) and the cosmic general ledger (cosmic consciousness) while the imbalance in the gratification deprivation complex always remain unbalanced when it comes to individual general ledger , advocating the evolutionary theory of individual consciousness to cosmic consciousness. It is to be noted that it is not the question of knowing all the transactions but the wisdom of what lies behind the debit- Credit, Debit-Debit, or Credit-Credit principle that is most important for assimilation.** Look at this simple thesis to prove the point stated in the foregoing. Thesis: 'The World Has A Beginning In Time, And In Space It Is Also Enclosed In Boundaries.' Proof: 'For If One Assumes That The World Has No Beginning In Time, Then Up To Every Given Moment In Time An Eternity Is Elapsed, And Hence An Infinite Series Of States Of Things In The World, Each Following Another, Has Passed Away. But Now The Infinity Of A Series Consists Precisely In The Fact That It Can Never Be Completed Through A Successive Synthesis. Therefore An Infinitely Elapsed World-Series Is Impossible, So A Beginning Of The World Is A Necessary Condition Of Its Existence, Which Was The First Point To Be Proved (Kant, Critique Of Pure Reason, Tr. By R. Guyer And A. W. Wood (Cambridge: Cambridge University Press, 1998): B 454, P. 470). As Quoted In The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO"The Cyclical Hypothesis, So Heavily Criticized By Nietzsche (VP II 325 And 334), Arises In This Way." In Fact, Nietzsche Was Not Criticizing The Cyclical Hypothesis But Only The Particular Form Of That Hypothesis Presented In Vogt's Work. All Of Nietzsche's Texts Without Exception Speak Of The Eternal Return As The Repetitio Of The Same Events Within A Cycle Which Repeats Itself

Eternally. If Deleuze's Interpretation Holds That The Eternal Return Is Not A Circle, Then What Is It? A Wheel Moving Centrifugally, Operating A "Creative Selection," "Nietzsche's Secret Is That The Eternal Return Is Selective" Says Deleuze: The Eternal Return Produces Becoming-Active. It Is Sufficient To Relate The Will To Nothingness To The Eternal Return In Order To Realize That Reactive Forces Do Not Return. However Far They Go, However Deep The Becoming-Reactive Of Forces, Reactive Forces Will Not Return. The Small, Petty, Reactive Man Will Not Return. Affirmation Alone Returns, This That Can Be Affirmed Alone Returns, Joy Alone Returns. Everything That Can Be Denied, Everything That Is Negation, Is Expelled Due To The Very Movement Of The Eternal Return. We Were Entitled To Dread That The Combinations Of Nihilism And Reactivity Would Eternally Return Too. The Eternal Return Must Be Compared To A Wheel; Yet, The Movement Of The Wheel Is Endowed With Centrifugal Powers That Drive Away The Entire Negative. Because Being Imposes Itself On Becoming, It Expels From Itself Everything That Contradicts Affirmation, All Forms Of Nihilism And Reactivity: Bad Conscience, Ressentiment..., We Shall Witness Them Only Once. [...] The Eternal Return Is The Repetition, But The Repetition That Selects, The Repetition That Saves. Here Is The Marvelous Secret Of A Selective And Liberating Repetition. There Is No Need To Remind The Reader That Neither The Image Of A Centrifugal Movement Nor The Concept Of A Negativity-Rejecting Repetition Appears Anywhere In Nietzsche's Writings, And Indeed Deleuze Does Not Refer To Any Text In Support Of This Interpretation. Further, One Could Highlight That Nietzsche Never Formulates The Opposition Between Active And Reactive Forces, Which Constitutes The Broader Framework Of Deleuze's Interpretation. For Some Years, Marco Brusotti Has Called Attention To The Fact That Deleuze Introduced A Dualism That Does Not Exist In Nietzsche's Writings. To Be Sure, The German Philosopher Describes A Certain Number Of "Reactive" Phenomena (For Example, In The Second Essay Of The Genealogy Of Morality, § 11, He Talks About

"Reactive Affects" [Reaktive Affekte], "Reactive Feelings" [Reaktive Gefühlen], Reactive Men [Reaktive Menschen]); But These Are Nonetheless The Result Of Complex Ensembles Of Configurations Of Centers Of Forces That Remain In Themselves Active. Neither The Word Nor The Concept Of "Reactive Forces" Ever Appears In Nietzsche's Philosophy. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO Some Parts Are Deleted Due To Spatial Constraints. In the midst of thoughts about the eternal return we find at least two other thematic axes On the one hand, the view of the **world as a constant flux of forces without any goal, law, or rules of becoming. A chaos sive natura de-divinized and de-anthropomorphized which constitutes the "ontological substratum**" of the whole of Nietzsche's reflexions On the other hand, an ensemble of fragments of an anthropologico-sociological character, designing a path of liberation leading to the creation of superior individuals by way of a profound transformation of their instinctual structure. This transformation must be achieved by a practice of solitude and internal struggle towards the liberation from the ancient representations of the world and from the incorporated herd values. For an analysis of these thematic perspectives, see Paolo D'Iorio, La linea e il circolo. Cosmologia e filosofia dell'eterno ritorno in Nietzsche (Genova: Pantograf, 1995): 233-322. With most of the parameters quantified, with assumptions there must be picturesque view of the self and non self, space and time, science and philosophy with the concatenational consummations. After The Chapter "On Redemption," Where Zarathustra Dares Not Expose His Doctrine, The Eternal Return Begins To Be Enunciated In Part Three Of The Work. In The First Place, It Is The Dwarf Who Formulates It In The Chapter "On The Vision And The Riddle." Facing The "Gate Of The Instant" Which Symbolizes The Two Infinities That Stretch Towards The Past And The Future, The Dwarf Whispers: "All Truth Is Crooked, Time Itself Is A Circle" The Dwarf Represents The Spirit Of Gravity, And He Embodies The Herd Morality, "The Belittling Virtue" Which Is The Title Of Another Chapter From Part III. The Dwarf Can

Endure The Eternal Return Without Great Difficulties Because He Has No Aspirations; Unlike Zarathustra He Does Not Wish To Climb The Mountains That Symbolize Elevation And Solitude. In Two Unpublished Notes, From The Summer And The Fall Of 1883, Nietzsche Writes: The Doctrine Is At First Favored By The Rabble, Before It Gets To The Superior Men. The Doctrine Of Recurrence Will First Smile To The Rabble, Which Is Cold And Without Any Strong Internal Need. It Is The Most Ordinary Of Life Instincts, Which Gives Its Agreement First. Hence, The Content Of The Doctrine Is The Same, But Whereas The Dwarf Can Endure It (Because He Interprets It According To The Pessimistic Tradition For Which "Nothing Is New Under The Sun"), Zarathustra, Who Is The "Advocate Of Life" Regards The Eternal Return As The Strongest Objection To Existence, And As The Rest Of The Dream Suggests, He Does Not Yet Succeed In Accepting It After The Vision At The Gate Of The Instant, The Chapter Is Brought To An End By The Enigma Of The Shepherd. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO .How does the memory affect the GD complex or for that matter the gratification deprivation of an individual or a collective identity? Non-Markovian local-in-time master equations give a relatively simple way to describe the dynamics of open quantum systems with memory effects. Despite their simple form, there are still many misunderstandings related to the physical applicability and interpretation of these equations. Here, we clarify these issues both in the cases of quantum and classical master equations. Further introduction of the concept of a classical non-Markov chain signified through negative jump rates in the chain configuration. **E-M Laine et al 2012 J. Phys. B: At. Mol. Opt. Phys. 45 154004 doi:10.1088/0953-4075/45/15/154004 Local-in-time master equations with memory effects: applicability and interpretation.** Supposing That There Were Indeed An "Energy Of Contraction" Constant In All Centers Of Force Of The Universe, It Remains To Be Explained Where Any Difference Would Ever Originate. It Would Be Necessary For The Whole To Dissolve Into An Infinite Number Of

Perfectly Identical Existential Rings And Spheres, And We Would Therefore Behold Innumerable And Perfectly Identical Worlds COEXISTING [Nietzsche Underlines This Word Twice] Alongside Each Other. Is It Necessary For Me To Admit This? Is It Necessary To Posit An Eternal Coexistence On Top Of The Eternal Succession Of Identical Worlds? The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. Science is usually portrayed as dehumanizing. Brave New World epitomizes this fear. "The more we understand the world, the more it seems completely pointless" (Steven Weinberg). Certainly science can seem chilling when conceived in the abstract as a metaphysical world-picture. We may seem to find ourselves living in a universe with all the human meaning stripped out: Participants in a soulless dance of molecules, or harmonics of pointlessly waggling superstrings and their braneworld cousins. Nature seems loveless and indifferent to our lives. What rights have we to be happy? **(Brave New World: Aldous Huxley).** Religion can afford to be dogmatic but science cannot. Look at this passage which counters both scientists and philosophers alike: However, Hartmann Objects That The Regressive Movement Postulated By Schopenhauer Is Possible Only In Thought: It Remains Nothing More Than An "Ideal Postulate" With No Real Object And Which "Does Not Teach Us Anything About The Real Process Of The World That Unfolds In A Movement Contrary To This Backwards Movement Of Thought" (Hartmann, Philosophie Des Unbewussten, Third Edition (1871): 772). Hartmann Affirms That If Unlike Schopenhauer One Admits The Reality Of Time And Of The World Process, One Must Also Admit That The Process Must Be Limited In The Past And Therefore That There Must Be An Absolute Beginning. In Hartmann's Mind, Failure To Do So Would Result In Positing The Contradictory Concept Of An Accomplished Infinity: "The Infinity That From The Point Of View Of Regressive Thinking, Remains An Ideal Postulate, Which No Reality May Correspond To, Must, For The World, Whose Process Is, On The Contrary, A Progressive Movement, Open Up To A Determinate Result; And Here The Contradiction Comes To Light" (Hartmann

(1871): 772). What Really "Comes To Light" In This Passage Is The Fact That Hartmann Does Not Provide A Demonstration But A Petitio Principii. Indeed, The Concept Of The World Process Analytically Contains The Concept Of A Beginning Of The World. In All Rigors, It Is Therefore Impossible To Demonstrate These Concepts With Reference To Each Other. Secondly, Hartmann's View That One Is Bound To Accept The Reality Of The World Process Even If One Rejects The Ideality Of Schopenhauer's Time Is Mistaken. Hartmann Believes That If Time Is Real There Must Be A World Process With Both An Absolute Beginning And An Absolute End. Without Any Justification, Hartmann Jumps From Schopenhauer's Negated Time To Oriented Time. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. In the last years several theoretical papers discussed if time can be an emergent property deriving from ) quantum correlations. Here, to provide an insight into how this phenomenon can occur, authors present an experiment that illustrates Page and Wootters' mechanism of "static" time, and Gambini et al. subsequent refinements. A static, entangled state between a clock system and the rest of the universe is perceived as evolving by internal observers that test the correlations between the two subsystems. They implement this mechanism using an entangled state of the polarization of two photons, one of which is used as a clock to gauge the evolution of the second: an "internal" observer that becomes correlated with the clock photon sees the other system evolve, while an "external" observer that only observes global properties of the two photons can prove it is static.Phys. Rev. A 89, 052122 (2014)DOI:10.1103/PhysRevA.89.052122 arXiv: 1310.4691 [quant-ph] **Time from quantum entanglement: an experimental illustration Ekaterina Moreva.** Time is an entanglement .When the new ideas of quantum mechanics spread through science like wildfire in the first half of the 20th century, one of the first things physicists did was to apply them to gravity and general relativity. The results were not pretty. It immediately became clear that these two foundations of modern physics were entirely

incompatible. When physicists attempted to meld the approaches, the resulting equations were bedeviled with infinities making it impossible to make sense of the results. Then in the mid-1960s, there was a breakthrough. The physicists John Wheeler and Bryce DeWitt successfully combined the previously incompatible ideas in a key result that has since become known as the Wheeler-DeWitt equation. This is important because it avoids the troublesome infinites—a huge advance. But it didn't take physicists long to realise that while the Wheeler-DeWitt equation solved one significant problem, it introduced another. The new problem was that time played no role in this equation. In effect, it says that nothing ever happens in the universe, a prediction that is clearly at odds with the observational evidence. This conundrum, which physicists call 'the problem of time', has proved to be thorn in flesh of modern physicists, who have tried to ignore it but with little success. Then in 1983, the theorists Don Page and William Wootters came up with a novel solution based on the quantum phenomenon of entanglement. This is the exotic property in which two quantum particles share the same existence, even though they are physically separated. Entanglement is a deep and powerful link and Page and Wootters showed how it can be used (e) to measure time. Their idea was that the way pair of entangled particles evolves is a kind of clock that can be used to measure change. But the results depend on how the observation is made. One way to do this is to compare the change in the entangled particles with an external clock that is entirely independent of the universe. This is equivalent to god-like observer outside the universe measuring the evolution of the particles using an external clock. In this case, Page and Wootters showed that the particles would appear entirely unchanging—that time would not exist in this scenario. But there is another way to do it that gives a different result. This is for an observer inside the universe to compare the evolution of the particles with the rest of the universe. In this case, the internal observer would see a change and this difference in the evolution of entangled particles compared with everything else is an important a measure of time. This is

an elegant and powerful idea. It suggests that time is an emergent phenomenon that comes about because of the nature of entanglement. And it exists only for observers inside the universe. **Time is an entanglement** (Wikipedia). Danish religious philosopher Soren Kierkegaard carried out a systematic critique of the pretensions of reason and an abstract rationalism which he believed that the modern age was nurturing. Condemning reflection as a "danger" that ensnares people in logical delays and machinations, Kierkegaard compared it to a prison. Reflection is for him a form of captivity, a bondage which "can only be broken by [passionate] religious inwardness" (1978: 81). **Reflection seduces individuals into thinking its possibilities are "much more magnificent than a paltry decision"** (1978: 82). It leads them to act **"on principle,"** to dwell on the deliberation of the context of their actions and the calculation of their worth or outcome. Kierkegaard argues that this drives away feeling, inspiration, and spontaneity, all of which are crucial for true inner being and a vital relation to God. For Kierkegaard, as Nietzsche would later agree, genuine inner being (and culture) is characterized by the tautness and tension of the soul which characterizes passionate existence. But the "coiled springs of life relationships ... lose their resilience" in reflection (1978: 78) and "everything becomes meaningless externality, devoid of [internal] character" (1978: 62). Kierkegaard thus contributes to the development of an irrationalist tradition that has echoes in some later postmodern thought. Kierkegaard might well have agreed with his contemporary Fyodor Dostoyevsky, who wrote: "An intelligent [reflective] man cannot seriously become anything ... excessive consciousness is a disease" (1974: 3, 5). In an age overtaken by rules and regulations, genuine action -- which Kierkegaard assumes to be subjective and spontaneous -- is frustrated at every turn. Complaining that we are too "sober and serious" (1978: 71) even at banquets, Kierkegaard bemoans the fact that even suicides are premeditated (1978: 68)! "That a person stands or falls on his actions is becoming obsolete; instead, everybody sits around and does a brilliant job of bungling through with the aid of some reflection

and also by declaring that they all know very well what has to be done" (1978: 73). Thus, it is passion, not reflection, that guarantees "a decent modesty between man and man [and] prevents crude aggressiveness" (1978: 62). "Take away the passion and the propriety also disappears" (1978: 64). The ambiguity in the word "passion" may cause some confusion here. To say that the age and its individuals are "passionless" is not to say there are no emotions whatsoever, but rather that there is no true spiritual inwardness and depth, no intensively motivated action and commitment. It suggests that passion exists only in a simulated, pseudo-form, "the rebirth of passion" through "talkativeness" (1978: 64). "Chattering" for Kierkegaard gets in the way of "essential speaking" and merely "reflects" inconsequential events (1978: 89-99). Hence, in "the present age," emotions -- which in fact are all too pronounced -- have been transformed into negative forces. Anticipating Nietzsche's genealogy of the "slave revolt" in morality, Kierkegaard claims that the "enthusiasm" of the prior age of Revolution, a "positively unifying principle," has become a vicious "envy," a "negatively unifying principle" (1978: 81), a leveling force in its own right insofar as those lacking in talent and resources want to tear down those who have them. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS By Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html.** In November 2011, the United States government argued before the US Supreme Court that it wants to continue utilizing GPS tracking of individuals without first seeking a warrant. In response, Justice Stephen Breyer questioned what this means for a democratic society by referencing Nineteen Eighty-Four. Justice Breyer asked, "If you win this case, then there is nothing to prevent the police or the government from monitoring 24 hours a day the public movement of every citizen of the United States. So if you win, you suddenly produce what sounds like 1984... "[65]. What is to be said in unmistakable terms is that when it comes to national security and innocent citizens there is question of compromise even it be an infringement and encroachment of so called

privacy which has been made a word to subterfuge the activities that are far said to be fair. Quantum theory has provoked intense discussions about its interpretation since its pioneer days. One of the few scientists who have been continuously engaged in this development from both physical and philosophical perspectives is Carl Friedrich von Weizsäcker. The questions he posed were and are inspiring for many, including the authors of this contribution. Weizsäcker developed Bohr's view of quantum theory as a theory of knowledge

## ABSOLUTE STABILITY

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range (0 W 9 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9 - 9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2 - 2y_9\frac{(E_9 + n_9^2 y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right) - 4y_9(E_9 - n_9)$$ 0

And has the solution

$$\frac{X_9}{M_9^2} > \frac{E_9}{(E_9 - 9)}[y_9(E_9 + n_9^2 y_9) + [y_9^2 (E_9 + n_9^2 y_9)^2 + 4y_9(E_9 - 9)(E_9 - n_9)]^{9/2}]$$

Given $E_9 > n_9$ this lower limit to $X_9 / M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9 / M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$ in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9 < n_9$ $(ie., 2A_9 M_9^2 < X_9 - Y_9)$ then (A2.9)and (A2.8)show that $Q_9$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2} > 2E_9[y_9 E_9 - (y_9 E_9 (9 + y_9)(E_9 - n_9)]^{9/2}]$$

And

$$E_9 > n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2 /2M_9 (X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$ and $X_9$ and $Y_9$ by $W_9^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9$ *and* separate (A3.1)into real and imaginary parts

$$Z_9^2 + (M_9 + V_9) Z_9 \sin(Z_9 T_9) - (U_9 M_9 + Y_9) + (M_9(U_9 + V_9) - X_9)$$

$$\mathrm{X}\cos(Z_9 T_9) - 2M_9 V_9 \cos^2(Z_9 T_9) + M_9 V_9 = 0$$

$$(M_9 + U_9) Z_9 - (M_9 + V_9) Z_9 \cos(Z_9 T_9) + (M_9(U_9 + V_9) - X)$$

$$\mathrm{X}\sin((Z_9 T_9) - M_9 V_9 \sin(Z_9 T_9) = 0$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 = 1 - \cos Z_9 T_9$and $A_9 = U_9 - V_9$we obtain

$$4 M_9^2 V_9^2 W_9^4 + 2 M_9 V_9 W_9^3[(M_9 - V_9)^2 + A_9(M_9 + V_9) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2[M_9 - V_9)^2(2 A_9 M_9 + Y_9 - X_9) + (2A_9 M_9 - X_9)^2 - A_9(M_9 - V_9)(2A_9 M_9 - X_9)$$

$$+ M_9 V_9 (A_9^2 + 4(X_9 + Y_9))]$$

$$+ W_9[A_9^2 + 2X_9)(2 A_9 M_9 - X_9) + 2A_9 Y_9(M_9 + V_9) - M_9 A_9^3] + A_9^2(X_9 M_9 - Y_9) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9(2, U_9 V_9) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting cos $(Z_9 T_9) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MDDULE NUMBERED SEVEN**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability

criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

**Morphological And Epistemological And Syllogistic Models For Which Nunny's Model Could Be Applied And Has Been Applied:**

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production**

**Charge Asymmetry, $\Delta Y \equiv Yt - Yt^{-}$, Reflects The Asymmetry In $Tt^{-}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles

ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic

quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation

between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to

nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of

Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first

impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of

another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease. Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation

of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces

always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the

thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of

thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity

and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI

Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs” Outcomes—

And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities

Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \quad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

## DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t-T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\big(R(t)\big)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and haimovici model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

**TIME DELAY**

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Limitations of Predator prey analysis in immense. When we realize that world itself is Debit-Credit world and sometime credit-Credit world (See an earlier paper entitled the same nomeneclature.

Opposite limit of absolute stability. The result is a specific The Doctrine of Vibration awareness of an object of the form: 'here this is such.' While individual appearances do not lose their separate identity {svarupabheda) when they rest on a common basis, even so the particular object which appears according to its own characteristics (svalak$ana) is an individual reality in its own right. It is a different kind of appearance characterised by its association with the appearance of the specific location and time in which it is made manifest. The form of our experience is thus 'I now see this here'. But when we perceive each particular constituent appearance separately, each assumes a separate fixed function. Abhinava cites the following colourful example to illustrate how the various combinations of appearances account for the variety of experience: Thus even though the appearance of the beloved may manifest externally, it is as if far away in the absence of another appearance, namely, that of 'embracing'. So when the [appearing of the beloved] is associated with another appearance [namely that of 'far away'] the power (arthakriya) it formerly had of giving pleasure appears as its contrary. The form our experience assumes depends, not only on the nature of the object perceived, but also on personal factors entirely peculiar to ourselves. This theory explains this in two ways. In one sense, the object remains the same, but one or other of its constituent appearances

comes to the fore according to the inclinations of the perceiver. From another point of view, we can say that the perceived object is different for each perceiver according to the difference in the prominent appearance manifest to him. Abhinava, citing as an example a golden jar, illustrates how the same object appears differently to different perceivers according to the use they wish to make of it and to their state of mind: When a person who is depressed and feels that there is nothing [of value for him in the world] sees the jar, he merely perceives the appearance 'exists' [in the form of the awareness that] 'it is\ He is not conscious of any other [of its constituent appearances] at all. An individual who desires to fetch water [perceives] the appearance 'jar'. The man who simply wants something that can be taken somewhere and then brought back [perceives] the appearance 'thing'. The man who desires money [perceives] the appearance 'gold'. The man who desires a pleasing object [perceives] the appearance 'brightness' while he who wants something solid sees the appearance 'hardness'. These 'atomic events' or appearances emerge from the pure subject's consciousness and combine together to form a total event at each moment. Integral Monism of Kashmiri Saivism Daily life (vyavahara) goes on by virtue of this ever renewed flux of appearances. They are connected together and work towards a single unified experience because they appear within the field of consciousness of the universal subject. The aggregate of appearances arises in the [supreme] subject as do [sprouts in] a rice field. Even though each sprout germinates from its own seed, they are perceived as a collective whole. Appearances rest in this way within the universal subject. 'External- ity' is itself another appearance; it arises from a distinction between appearances and the individual subject. So, although all manifestation always occurs within the subject, it appears to be external due to the power of Maya which separates the individual subject from his object. This split must occur for daily life to be possible. Only externally manifest appearances can perform their functions; when they are merged within the subject and at one with him, they cannot do so. Daily life proceeds on the basis of the operation and withdrawal of the conditions necessary for fruitful action to be possible. Appearance in this sense represents the actualisation of a potential hidden in consciousness made possible by virtue of its dynamic, Spanda

nature which is both the flow from inner to outer and back as well as the power that impels it. The emergence from, and submergence into, pure consciousness of each individual appearance is a particular pulsation (visesaspanda) of differentiated awareness. Together these individual pulsations constitute the universal pulse (samanyaspandd) of cosmic creation and destruction. Thus, every single thing in this way forms a part of the radiant vibration {sphuratta, sphurana) of the light of absolute consciousness. Light and Awareness: The Two Aspects of Consciousness Absolute consciousness understood as the unchanging ontological ground of all appearing is termed 'Prakasa'. As the creative awareness of its own Being, the absolute is called ' Vimarsa '. Prakasa and Vimarsa — the Divine Light of consciousness and the reflective awareness this Light has of its own nature — together constitute the all-embracing fullness (purnata) of consciousness. The Recognition (pratyabhijfia) school of Kashmiri Saivism develops this concept of the absolute which finds its fullest expression in Utpaladeva's Stanzas on the Recognition of God. Even though neither of these two key terms appear in the Stanzas on Vibration or the Aphorisms of Siva, they recur frequently in their commentaries. Thus, although the original formulation of the Doctrine of Vibration differs from the theology of Recognition in this respect, it was extended in the course of its development to accommodate this concept of the absolute as well. This was possible, and quite justified, insofar as the absolute understood in Pratyabhijfia terms does not, as we shall see, differ essentially from that of the Spanda school. We can, as Kashmiri Saivites themselves have done, explain one in terms of the other. The Doctrine of Vibration Prakasa: The Light of Consciousness Prakasa is the pure luminosity' (bhdna) or 'self-showing' that constitutes the essence and ultimate identity (atman) of phenomena. That things appear at all is due to the light of consciousness, and their appearing (avabhasana) is itself this Light which bestows on all things their evident, manifest nature. Established in the light of consciousness everything appears there according to its own specific nature (svabhava). Anything that supposedly does not rest in this Light is as unreal as a sky-flower. 3 Thus, according to Rajanaka Rama, unlike the light of the sun, or any other light, this Light not only makes all things apparent, it is also their ultimate source. 4 Full of its

divine vibration the Light makes all things manifest and withdraws them into itself. This supra- temporal activity characterises it most specifically; devoid of it, it would be no better than an inert physical phenomenon. At the same time, this light is the conjunction (slesa) or oneness (aikdtmya) of its countless manifest forms, 6 and the collective whole (Sampina'ana) of all the categories of existence. The universe is nothing but the shining of the Light within itself. It is the radiant vibration (sphuratta) of this Light, the state (avast hand) in which consciousness becomes manifest. Although the Light shines as all things at all times and hence also makes their diversity manifest, 9 penetrating each object individually as well as collectively, it is not totally 'merged' (magna) or identified with the object so as to suffer any division within itself. Our experience of any object is of the form: T see this': it is not itself an object, but the manifest form the object assumes as a luminous principle of experience. The Light is ever revealed and can never be obscured; objectivity can never cast a shadow on the light of consciousness. The Stanzas on Vibration declare: That in which all this creation is established and from whence it arises is nowhere obstructed because it is unconditioned by [its very] nature. This Light is the highest reality (Paramartha). lity with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**DELAY DUE TO CONFEDERATION DEBIT TYPE MATURATION CREDIT TYPE CONSOLIDATION**

**MATURATION**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \quad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9} + \frac{N_9 dG_9}{dR_9} - \frac{dB_9}{dR_9} > \frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9} + \frac{dM_9}{dR_9}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened**

**eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production Charge Asymmetry, $\Delta Y \equiv Yt - Y\bar{t}$, Reflects The Asymmetry In $Tt^{-}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton

particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results

in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely

dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact**

**most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the**

**thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is

nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the

world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease.

Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition

and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the

second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular

chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon

be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining

Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been

Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing

Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

**ABSOLUTE STABILITY**

$$L_9 = \frac{N_9 G_9}{2M_9} \frac{d(9nF_9)}{dR_9} (\geq 0) \quad (9a)$$

$$I_9 = -\frac{N_9 G_9}{2M_9} \frac{d(9nM_9}{dR_9} (\geq 0) \quad (9b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9 > L_9 - I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_9 > L_9 - I_9,$ when $P_9^2/4 > I$ (8a)

$P_9 > L_9 - I_9 + \frac{(I_9 - P_9^2}{(L_9 - I)}$ when $P_9^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to

reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

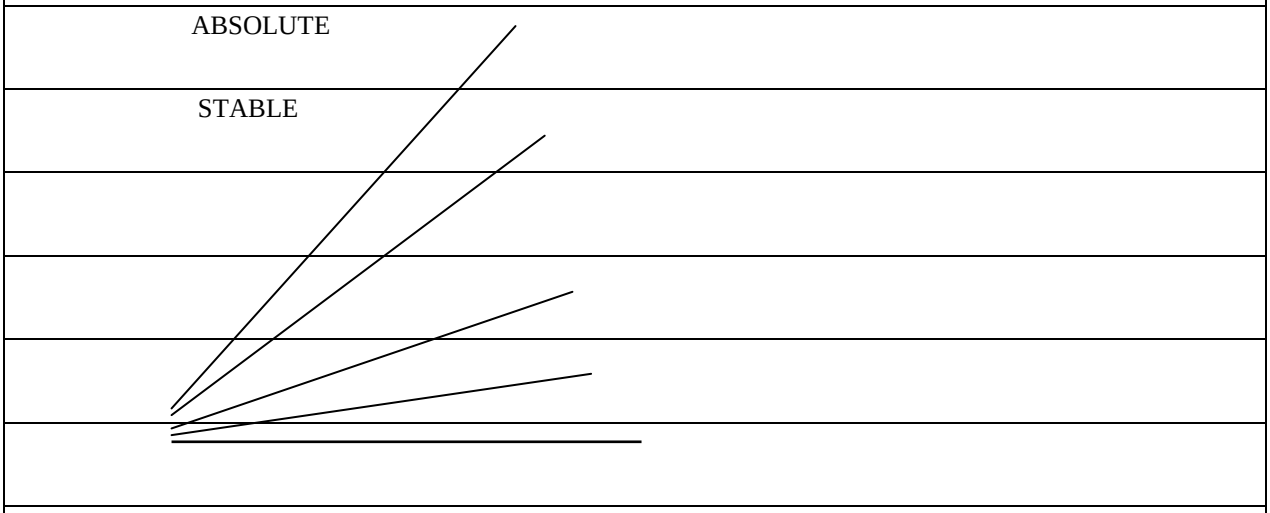



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori

reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

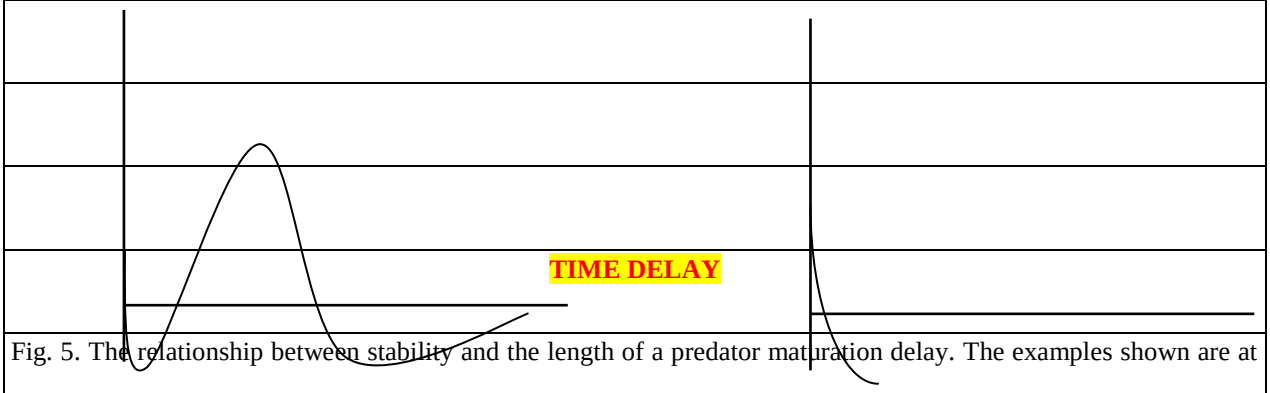



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

**DELAY IN BOTH PREDATOR AND PREY VIS A VIS DEBIT TYPES AND CREDIT TYPES**

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t - T_9)F_9\big(R_9(t - T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t - T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) and the effective prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting

$dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

Systems to which model has been applied also include: **Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI:

http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors.

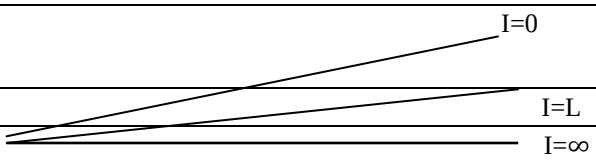
Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, ,**

**Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)**

TOTAL

ABSOLUTELY STABLE

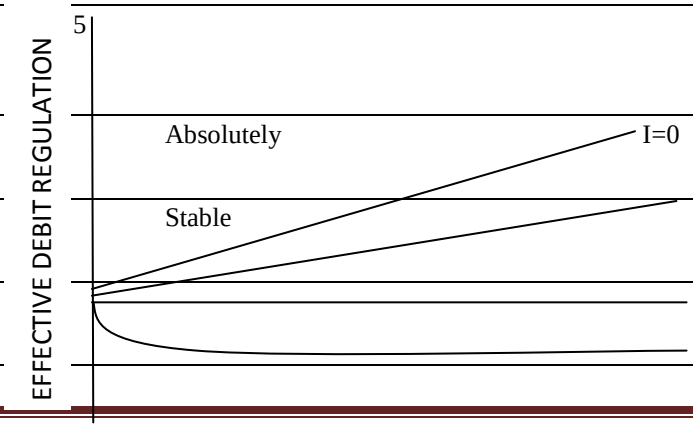



DELAYED DEBIT EFFECT

FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.



-50 10

**DELAYED CREDIT EFFECT (LACK OF BALANCE) CREDIT NON CUMSUMMATION:"**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

The following treatise bears ample testimony and infallible observatory to the fact that how Nunny's model can be applied to various cosmological, general relativity and economic problems.

## ONTOLOGICAL AND EPISTEMOLOGICAL FEATURES OF THE STUDY: OF NUNNY's MODEL

When they are known to be as they are, in the fullest sense of their existence, their phenomenal nature disappears leaving the ground of Being naked and accessible. This approach was validated by a critique of experience. The Vedanta established that space, time and the other primary categories of our daily experience can have no absolute existence. It was therefore necessary to make a distinction between relative truths — that accepted by the precritical common man — and an absolute truth discovered at a higher level of consciousness. The Saiva absolutist 17 rejects any theory that maintains that the universe is less than real. From his point of view a doctrine of two truths, one absolute and the other relative, endangers the very foundation of monism. The Kashmiri Saiva approach is integral: everything is given a place in the economy of the whole. It is equally wrong to say that reality is either one or diverse. Those who do so fail to grasp the true nature of things which is neither as well as both. "We do not" says Abhinavagupta, "base our contention that [reality] is one because of the contradictions inherent in saying that it is dual. It is your approach {paksa) that accepts this [method] While, if [duality and oneness] were in fact [to contradict each other], they would clearly be two [distinct realities]." The Vedantin, who maintains that non-duality is the true nature of the absolute by rejecting duality as only provisionally real, is ultimately landed in a dualism between the real and illusory by the foolishness of his own excessive sophistry (vacafadurvidya). Oneness is better understood as the coextensive unity (ekarasa) of both duality and unity. They are equally expressions of the

absolute. Gopinath Kaviraj says: According to Sarikara, Brahman is truth and Maya is inexplicable (anirvacaniya). Hence the [Advaitin's] endeavour to demonstrate the superiority of Advaita philosophy is turned against his own system. It tarnishes the picture of its philosophical perfection and profundity. He cannot accept Maya to be a reality, therefore his non-dualism is exclusive. The whole system is based on renunciation and elimination and thus is not all-embracing.... By accepting Maya to be Brahman (brahmamayi), eternal (nitya) and real (satyarupa), Brahman and Maya [in the Tantra] become one and coextensive. **The Doctrine of Vibration: An Analysis of the Doctrines and Practices of Kashmir Shaivism Mark S. G. Dyczkowski. Concept of being from Shaivite point of view:** Solutional behaviour and stability analysis of such systems is of paramount importance, notwithstanding the assumptions made which are not farfetched, in that it throws light on the fact whether they could be candidates for yet unsolved problems, providing Autheticative determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication for the wide amplitudinal ramificatory usage of the thesis propounded. The world is now the other person's world, a foreign world that no longer comes from the self, but from the other. The other person is a "threat to the order and arrangement of your whole world...Your world is suddenly haunted by the Other's values, over which you have no control"(Satre: being and Nothingness) Sartre states that many relationships are created by people's attraction not to another person, but rather how that person makes them feel about themselves by how they look at them. This is a state of emotional alienation whereby a person avoids experiencing **their subjectivity by identifying themselves with "the look" of the other**. The consequence is **conflict**. In order to maintain the person's own being, the person must control the other, but must also control the **freedom of the other "as freedom**". These relationships are a profound manifestation of "bad faith" as for-itself is replaced with the other's freedom. The purpose of either participant is not to exist, but to maintain the other participant's looking at them. This system is often mistakenly called "love", but it is, in fact, nothing more than emotional alienation and denial of freedom. In

essence we are always simulating and selling ourselves at the cost of our freedom. Both as a script writer and participant consciousness, human beings experienceâ complete subjective feelings and sometimes absolute subjective feelings, which if remains unchecked might lead to pathological ramifications. (Italics mine) Perspicuous forbearance, Sophisticated seasoning, Partcipational observation, Background actuality, Existential worldliness, Existential strife, Predicational anteriorities, Character consonance, Ontological consonance, Primordial exactitude, Phenomenological correlates, Accolytish representation, Atrophied asseveration. Anamensial alienisms, Anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orinetatilities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state form the bastion, pillar, post, stylobate and sentinel of the classification scheme and doxa. appropriate and diagnostic, differentiating-discriminating, distinctive-distinguishing, emblematic- especial, essentially exclusive, fixed, idiosyncratic-inborn, inbred, indicative and individual, individualistic-inherent, individualizing and ingrained, inherent, innate, local, marked, native, normal, original, particularly- peculiar, personal, private, proper, regular, representative-singular, special, specific, symbolic, symptomatic- unique these form the classification and stratification measure that is suitable notwithstanding the constancy in many a case, as in gravity, which depends on the mass and distance for two given objects**.** Like Kant and Bergson, Deleuze considers traditional notions of space and time as unifying forms imposed by the subject. He therefore concludes that pure difference is non-spatio-temporal; it is an idea, what Deleuze calls "the virtual". (The coinage refers to Proust's definition of what is constant in both the past and the present: "real without being actual, ideal without being abstract.") Proust, Le Temps Retrouvé, ch. III: see the fourth line from the bottom of this page, or, in English translation, the thirteenth paragraph here: "I began to discover the cause by comparing those varying happy impressions which had the common quality of being felt simultaneously at the actual moment and at a distance in time, because of which common quality the noise of the spoon upon the plate, the unevenness of the paving-stones, the taste of the madeleine, imposed the past

upon the present and made me hesitate as to which time I was existing in. Of a truth, the being within me which sensed this impression, sensed what it had in common in former days and now, sensed its extra-temporal character, a being which only appeared when through the medium of the identity of present and past, it found itself in the only setting in which it could exist and enjoy the essence of things, that is, outside Time. [...] **Nothing but a moment of the past Much more perhaps; something which being common to the past and the present, is more essential than both. [...] a marvellous expedient of nature had caused a sensation to flash to me—sound of a spoon and of a hammer, uneven paving-stones**—simultaneously in the past which permitted my imagination to grasp it and in the present in which the shock to my senses caused by the noise had effected a contact between the dreams of the imagination and that of which they are habitually deprived, namely, the idea of existence—and thanks to that stratagem had permitted that being within me to secure, to isolate and to render static for the duration of a lightning flash that which it can never wholly grasp, a fraction of Time in its pure essence. When, with such a shudder of happiness, I heard the sound common, at once, to the spoon touching the plate, to the hammer striking the wheel, to the unevenness of the paving-stones in the courtyard of the Guermantes' mansion and the Baptistry of St. Mark's, it was because that being within me can only be nourished on the essence of things and finds in them alone its subsistence and its delight. It languishes in the observation by the senses of the present sterilised by the intelligence awaiting a future constructed by the will out of fragments of the past and the present from which it removes still more reality, keeping that only which serves the narrow human aim of utilitarian purposes. But let a sound, a scent already heard and breathed in the past be heard and breathed anew, simultaneously in the present and in the past, real without being actual, ideal without being abstract, then instantly the permanent and characteristic essence hidden in things is freed and our true being which has for long seemed dead but was not so in other ways awakes and revives, thanks to this celestial nourishment "While Deleuze's virtual ideas superficially resemble Plato's forms and Kant's ideas of pure reason, they are not originals or models, nor do they transcend possible experience; instead

they are the conditions of actual experience, the internal difference in itself. "The concept they [the conditions] form is identical to its object."[Desert Islands, p. 36.] A Deleuzean idea or concept of difference is therefore not a wraith-like abstraction of an experienced thing, it is a real system of differential relations that creates actual spaces, times, and sensations. [See "The Method of Dramatization" in Desert Islands, and "Actual and Virtual" in Dialogues. In fact it not as if the style and substance of science and philosophers have been entirely different. look at the conformality and consonance in the Gaussian Theory and Deleuze's statements on identity and the loss of it in the following passage…….That is, identity is a continuous process that finds its limit in exhaustion. As Deleuze will write, "…the logical relation of causality is inseparable from a physical process of signalling, without which it would not be translated into action. By 'signal' we mean a system with orders of disparate size, endowed with elements of dissymmetry; by 'sign' we mean what happens within such a system, what flashes across the intervals when a communication takes place between disparates. The sign is indeed an effect, but an effect with two aspects: in one of these it expresses, qua sign, the productive dissymmetry; in the other it tends to cancel it" (dr, 20). An actualized entity just is such a signal-sign system, remaining in perpetual communication with the virtual multiplicity it actualizes, constituting its own elements and identity. Such is Deleuze's account of what Badiou refers to as the operation of the "count-for-one"… an operation that is immanent to being itself and energetic in character. yet how are we to discern the local nature of ontological situations, the manner in which there is no global, overarching global situation in which all local situations are embedded as if parts of a whole? in part the above passage already responds to this question. if series must be brought into relation to communicate in order for causality to occur, and if the resonance of series is always a matter of chance, of a throw of the dice, like Lacan's objet a where no series enjoys primacy of model over copy, then all we have are local situations without a total situation. Unlike Heidegger's analysis of Dasein where Dasein is always being-in-the-world, Deleuze's simulacra or actualizations are always divergent, enjoying only local relations. Heidegger argues that in order to

understand the being of Dasein we must understand the manner in which Dasein is a being-in-the-world. By contrast, Deleuze's concept of multiplicity provides us with the principle of a new structuralism, of a new local ontology, that allows us to understand the immanent organization of a multiplicity without referring it to an embedding global space. as DeLanda so beautifully puts it in relation to the mathematician gauss, "…when gauss began to tap into these differential resources, a curved two-dimension surface was studied using the old Cartesian method: the surface was embedded in three-dimensional space complete with its own fixed set of axes; then, using those axes, coordinates would be assigned to every point of the surface; finally, the geometric links between points determining the form of the surface would be expressed as algebraic relations between the numbers. but gauss realized that the calculus, focusing as it does on infinitesimal points on the surface itself (that is, operating entirely with local information), allowed the study of the surface without any reference to a global embedding space. basically, gauss developed a method to implant the coordinate axes on the surface itself (that is, a method to implant the coordinate axes on the surface itself (that is, a method of 'coordinatizing' the surface) and, once points had been so translated into numbers, to use differential (not algebraic) equations to characterize their relations. as the mathematician and historian Morris Kline observes, by getting rid of the global embedding space and dealing with the surface through its own local properties, 'gauss advanced the totally new concept that a surface is a space in itself'" (intensive science and virtual philosophy, 11-12).according to the old Cartesian method, we can only outline the properties of a space by relativizing it to a global space in terms of which it is then mapped. by contrast, gauss is able to explore a space in terms of its intrinsic metric and organization as a local space, without referencing it to a whole of which it is conceptualized as a part. **The Image of Thought** "Image of thought "permeates both popular and philosophical discourse. According to this image, thinking naturally gravitates towards truth. Thought is divided easily into categories of **truth and error**. The model for thought comes from the educational institution, in which a master sets a problem and the pupil produces a solution which is either true or false. This image of the subject

supposes that there are different faculties, each of which ideally grasps the particular domain of reality to which it is most suited. In philosophy, this conception results in discourses predicated on the argument that "**Everybody knows..."** the truth of some basic idea. Descartes, for example, appeals to the idea that everyone can at least think and therefore exists. Deleuze points out that philosophy of this type attempts to eliminate all objective presuppositions while maintaining subjective ones. Deleuze maintains, with Artaud, that real thinking is one of the most difficult challenges there is. Thinking requires a confrontation with stupidity, the state of being formlessly human without engaging any real problems. One discovers that the real path to truth is through the production of sense: the creation of a texture for thought that relates it to its object. Sense is the membrane that relates thought to its other. Accordingly, learning is not the memorization of facts but the coordination of thought with a reality. "**As a result, 'learning' always takes place in and through the unconscious, thereby establishing the bond of a profound complicity between nature and mind".** Deleuze's alternate image of thought is based on difference, which creates a dynamism that traverses individual faculties and conceptions. This thought is fundamentally energetic and signifying: if it produces propositions, these are wholly secondary to its development. Philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; even tum tan tum;it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the vorticity of being merges with the positive use of the disjunctive synthesis, and

this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoin der for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for every hum of voices or a single silence for all the silences; a single vacuum for all the vacuousness consciousness should not be said without conjuring if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. The attributable to all states of bodies and states of affairs and the expressible of every proposition So univocity of consciousness means the identity of thé noematic attribute and that which is expressed linguistically and sense-fully. Univovity.We mean that it does not allow consciousness to be subsisting in aquasi state and but expresses in all pervading reality. (Reference: LOS: Deleuze) Deleuze sometimes spoke of getting up behind an author and "creating a monster", as if his reading methodology somehow consisted of a distortion of the philosopher's or author's thought. Certainly there is a case to be made for this with regard to the process of deterritorialization, where something is wrested from a territory and deterritorialized upon a new territory, like the animal paw that is deterritorialized from the earth and reterritorialize on the branch. however, it seems to me that the more interesting aspect of Deleuze's approach to other philosophers and art is not so much his "monstrous creations" (i seldom find them particularly monstrous), but rather their Gaussian or Riemannian style, where he explores them in terms of their own internal organization and metric, without reducing them to something alien such as history, society, biology (reductivism), or the signifier. What we find in Deleuze's approach to phenomena is a Gaussian technique. for instance, take Deleuze's books on cinema, his book on Francis bacon, or his study of sacher-masoch. In the first instance, Deleuze's carefully separates cinema from narrative and the signifier, studying it in terms of its specific organization pertaining to the production of images. in the case of bacon, Deleuze doesn't look for an underlying narrative or "meaning", but instead studies the manner in which bacon composes and

organizes his images and lines, both in terms of their production and actuality, so as to liberate a logic of sensation. finally, in approaching sacher-masoch, Deleuze allows Sacher-Masoch's novels to speak for themselves in terms of their desire and relation to pain, stalwartly refusing to reduce masochism to the complement of sadism. in each case, we have a local exploration of a "space" of multiplicities that is extremely precise. Sean Bowden the priority of events: **Deleuze's logic of sense sean Bowden, the priority of events: Deleuze's logic of sense, Edinburgh university press, 2011, 296pp., $40.00 (hbk), isbn 9780748643646 Reviewed by David Scott, Coppin state university and Gaussian spaces and multiplicities a note on Gaussian spaces and multiplicities posted by larval subjects under uncategorized.** What you "see" is not always what you "get." Many people mistakenly take their own visions literally without "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in the physical sense but in a psychological sense. The archetypes of the collective unconscious are arrayed behind the scenes of current worldwide conditions, of crisis and confusion. They mirror our own states back at us, whether we perceive them as such or interpret them plausibly or not. The noise of ordinary consciousness and beliefs drowns out the signal. Unconsciousness is the background of our ordinary awareness. Our organism is very much at the center of such effects. The organismic source is our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is (=) a meme-Scape. Pre-conceived concepts vie with (e&eb) structures, concepts with images. Like scientists who ignore (e) assumed truths, we leapfrog over (e&eb) our beliefs and personality deficits, claiming (eb) idiosyncratic imagination is (=) literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves (e) us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our

experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images. It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is (=) both. We inherit (e&eb) the structures that make our experience (=) what it is. But the structure itself is "empty," and each human culture "fills" it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about (e&eb) our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from

criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns,

nomenclature, and internal coherence of meaning and representation. n a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar some statements in Vedanta, it is not superfluous or redundant. It is evolutionary. Authetication determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication of the evolutionary process towards understanding is ever present in all, its thematic and discursive form. it is the relentless holding on to a theory this is despicable albeit the theory is not. Look at this paragraph. It is the 'I am the body/mind' belief that gives rise to the 'I am not (e) the world' belief. These two beliefs are (=) co-created."Consciousness projects (e&eb) the appearance of the mind, body and world by taking the shape of thinking, sensing and perceiving." "Attention is Consciousness with an object. When the object vanishes, attention simply remains what it always is, Consciousness.""There is (=) no purpose to (e) meditation. The purpose is already accomplished." "Everything that is experienced is experienced by, through, in and as Consciousness." "The seen cannot be separated from seeing and seeing cannot be separated from (e) Consciousness.""The Reality of any experience is not hidden in the appearance; it is expressed by the appearance." "Once we see that everything is Consciousness... Maya still dances, but it is a dance of love not seduction." There are some truly excellent sections

such as a long one taking, as a starting point, the following quotation from the painter Cézanne: "Everything vanishes, falls apart, doesn't it? Nature is always the same but nothing in her that appears to us lasts. Our art must render the thrill of her performance, along with her elements, the appearance of all her changes. It must give us a taste of her Eternity." Rupert begins: "That statement must be one of the clearest and most profound expressions of the nature and purpose of art in our era." And he goes on for some 14 pages to elaborate on this claim, examining the nature of the 'elements' Existence, appearance and Consciousness and their relationship

Some of the following concepts are modeled; some are not. For those which are not modelled kindly see one of the papers in the series. This note has the axiomatic predication to give a holistic view of the study. Here we talk of the characterstics of systems which satisfy the condition of cosmological constant or any other for that matter or variable or axiomatic predication. . There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit,

facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified

from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. Towards the end of circumvention of typing hundreds of superscripts and subscripts which would be a sardine squash and the concomitant operational difficulties, model is presented in piece meal of nine modules. Logarithms are to be taken in respect of those which incorporate more than one variable in bra-ket. Values of log (ab) and log (a+b) are readily available. Anti logarithm shall be taken at the end to predict, project, prognosticate the value of the variable. This is true for tensors, vectors and other variables too. Affirmational assertion, and explanation justification, statement of vindication, annotational commentary, and explicational glossary for each and every system changes and physical interpretation of results is one thing that is to be with earnest endeavour and feverish and febrile expediency. Editions never ever mean the same and identicality of parameters and this has been explicitly and unmistakably stated in the model apriori itself. Akin and analogous, cognate and concurrent, correspondent and congruent, comparable and complementary, synonymous and duple, tantamount and agnate, commensurate and correlative representation is only to highlight the importance and subterfuge the replication of work. It is my fervent solicitation to kindly bear with me for any lapses, notwithstanding the orchestrated efforts for a paper without any minor errors. Postulation predication, conclusive presumption, differential presuppositions, underscored decidedly axiomatic statement of the statement, equation form the bye word or the watch word in the aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement-argument, articulation, assertion, asseveration, choice of words, commentary and communication, declaration and definition, delivery-diction, elucidation-emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation tone and tenor of stratification. Now, how do find the reaction of systems to these singularities. You do the same thing a boss does for you. "Problematize" the events and see how you behave. I will resort to "pressure tactics". "intimidation of deriding report", or "cut in the increment" to make you undergo trials, travails and tribulations. I am happy to see if you improve your work; but may or may not be sad if you succumb to it and hang yourself! We do the same thing with systems. systems show conducive response, felicitous reciprocation or behave erratically with inner roil, eponymous radicality without and with blitzy conviction say like a solipsist nature of bellicose and blustering particles, or for that matter coruscation, trepidiational motion in fluid flows, or seemingly perfidious incendiaries in gormandizing fellow elementary particles, abnormal ebullitions, surcharges calumniations and unwarranted(you think so but the system idoes not!) unrighteous fulminations. Amount of emotion or affect that is appropriate or expected in an given form of interaction. Particular individuals and diffuse obligations (see c and d) are associated with affectivity, whereas contacts with many individuals (universalistic) in a bureaucracy may be devoid of emotion and characterized by affective neutrality. Affective neutrality may refer to self discipline and the deferment of gratification (eg Weber's spirit of capitalism). In contrast, affectivity may be associated with expressing emotions. Adams and Sydie also refer to affective neutrality being associated with ego control (p. 15). (Parsons: Wikipedia). So the point that is made here is "like we problematize the "events" to understand the human behaviour we have to "problematize" the events of systems to understand their behaviour. This statement is made in connection to the fact that there shall be creation or destruction of particles or complete obliteration of the system (blackhole evaporation) or obfuscation of results. Some systems are like "inside traders" they will not put signature atoll! How do you find they did it! Anyway, there are possibilities of a CIA finding out as they recently did! So we can do the same thing witn systems to. This is accentuation, corroboration, fortification, .fomentatory note to explain the various coefficients we have used in the model as also the dissipations called door. In the bank example we have clarified that various systems are individually conservative, and their conservativeness extends holisticallytoo.that one law is universal does not mean

there is complete adjudication of nonexistence of totality or global or holistic figure. Total always exists and "individual" systems always exists, if we donot bring Kant in to picture! For the time being let us not! Equations would become more ensorcelled and frenzied..... philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". In consciousness research, two rival sets of theories can be recognized: (A) Scientific material interpretations of consciousness are based on axioms that view consciousness in the context of highly advanced intentional processing of information in which subject-object relations evolve, and (B) humanistic interpretations of consciousness are based on axioms that view consciousness in the context of, say, "centered pulsations" that enable a conscious agent to act from his or her center of awareness. In this paper I will argue for the selection of axioms that favor humanistic interpretations of consciousness. The Pursuit of Autonomy Interdisciplinary Observations to Human Consciousness: Wautischer, Helmut (1993) the Pursuit of Autonomy. Interdisciplinary Observations to Human Consciousness: Social Neuroscience Bulletin, 6 (4). pp. 52-56.Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality; **CONSCIOUSNESS AND ITS UNIQUENESS June 12, 2012 at 8:40pm (See Deleuze Logic of sense, Wikipedia and Stanford encyclopedia for more details)** Atrocious contrivance and device, gratificational primogeniture, calamitous dodge and expedient, depraved, destructive, disastrous, execrable gambit and gimmick, great solace and succor, iniquitous, injurious, loathsome, low, maleficent, malevolent machination and maneuver, promethaleon of candor, frankness, honesty, honor, ingenuousness, innocence, openness, reality, sincerity, truthfulness, spiteful, stinking, ugly, unpleasant, unpropitious play and ploy, progenitor of jurisprudence and circumspection, racket and ruse, savvy, scam, stratagem, subterfuge, tactic and wile, forthrightness, honesty, truthfulness embodied and personified is the structural predisposition and dispensation of the complex. Nobody need have to take the Brahman-Anti Brahman agency to be a recalcitrant, repugnant, and refractory proposition. God gives and also takes away. That is the point made. ‘

And of course lesser gods, and God doth follow them and they follow God. That is why they are lesser gods. And wield the power on lesser mortals like us. When we mean gratification and deprivation we mean celestial Brahman Anti Brahman and Terrestial Brahman and Anti Brahman agencies producing such happiness or sadness by any means such as (blending- amalgamation, adulterant-adulteration, amalgam-amalgamation, blend- combination,

composite- compound, debasement-denaturant, fusion-hybrid, intermixture-reduction, amalgam- mixture, composite- compound, fusion-mishmash, aggregate- alloy, amalgamation-blend, commixture-composite, composition-compost, conglomerate-fusion, goulash-medley, mishmash-stew, blend-coalition, commixture-compound, federation-heating, integration-junction, melting-merger, synthesis-unification, assortment- combination, adulteration-alloy, assimilation-association, batter-blend, brewed combine, composite compound, concoction confection, conglomeration cross of money, orientation(dis), penury creation, murder, mayhem plunder, pillage, apocalypse, Armageddon in various forms and permutation and combination of Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness: with each action accounted and showing balance in the nature's general ledger as also actionlessness:e). All these are recorded in Nature's general ledger. Neuron DNA encapsulates all the actions passions, actions, interactions, transactions or the lack of it at various levels of individual, collective and cosmic levels. In celestial Brahman Anti Brahman ledger there is cessation of dualities. One more point to be made is if 5=3+2, then we can say 5 gormandizes 3 or 5 gobbles up 2. Quintessentially such a statement arises out of the difference of LHS and RHS term wise. It must be remembered despite all theoretical abstractions and generalizations be it proton mass or Higgs Boson, or topology of a space, it is notwithstanding the fundamentality of nature is just a number. Just everything in Physics and mathematics is just a number. Let us assume I buy rice in Bangalore (District) worth 5k.g. I can always divide 5k.g.s of rice in categories of Bangalore

(South), Bangalore (west), Bangalore (East) and Bangalore (South). On similar basis, there are billions of neutrons, protons and concomitant interactions which would satisfy the given numerical value. Example of gravity above makes things clear. Everything atleast is a function of time and someplace or some other coordinates, subordinates or superordinates. Mass of proton is 2.47 does not mean total mass of all protons is 2.47. It is this argument that is used in the classification scheme. By the very nature, self the unmonitoring but the witness consciousness agency identifies itself with something, be it a profession, an actor on the screen, when it does not have a identity. Such examples are legion, a Charlie Chaplin in "The Kid" or Raj kapoor in "Jagte Raho". These are people who just wanted to live and living without being put under surveillance every moment is the greatest part of their life. This is "Absolute Subjectivity" of Shiva. Such a state is always craving for identification for in pure consciousness all identities are lost. It is this "identitylessness" or having very minimal identity is what pervades most in India and this is due to expansion of individual consciousness. When that happens, be it education, or wealth, the "identity' becomes so well entrenched, they become the exact opposite of "pure consciousness"," impure consciousness" (Shakti). This is a state of "Absolute or relativistic objectivity". There was this case of a lady who went to such an extent of planning, she was sure that whatever she does must work out. When it did not, she lapsed in to "identity crisis". Such thing as this might happen in well developed countries where the competition is extremal and to be achieved at any cost. Realisation that one is not the doer but the witness consciousness is just like watching a movie in which one has acted. It is in this sense "advaita", "non duality' seems to be of cardinal and preeminent importance for the analytical treat ment of "identity' and "identity lessness", both happen in space time. "Multiverses" in this sense is also like a different film of an actor released at different theaters and it is in this sense we state that that every object in space time is a multiverse. You must have totally immersed in doing a "string theory" problem or watching a "film", oblivious of surroundings. Essentially creation (space time) also is what is happening in mind and honed in to you billion times. Identification of anything with a designation or signification is one that leads to

association within spacetime, which would only have pathological consequences and detrimental ramification as stated above or positive implications of "happiness" because of the identity itself. Search or 'show" has to go on! It gives a sense of consummation. a watching pot never boils; donot go to check nature; you are disturbing it; it will never show its true colours bhakti and Virakthi, gratification and deprivation, raga and dwesha is essential to live in space. It is the quantum interest thereof that produces energy to live .gratification and deprivation which form the bastion, pillar, post and stylobate and sentinel of all human actions depends on tams, rajas, and sattva of individual or insentient object. Verily tamas, rajas and satva are the characterstics of space itself, nay the behavioural pattern and attitudinal orientation. There are peaceful places like Vienna or Sringeri and there are violent places like others. Space is warped and woofed by the Akshara, the quantum information of primordial syllable Om. cosmic general ledger is the one vibration of which is brahman; quantum information possessions thereof form the sentinel of all transactions (Gargi to Yajnavalkya: my interpretation)adhocism, nonseqientiality of events, anti sequentiality of events, simultaneity of phenomenon, concurrent happenings , improbability of phenomenon happening itself are characterstics of brahman-anti brahman (Terrestial and celestial) are the cause of non predictability of happenings like in black hole, or in the areas of operation of brahman-anti brahman (Terrestial and celestial). brahman AntiBrahman actions, interactions, decisions, circumspection, adhocism, satisfies sensitive to initial conditions; topologically mixing; dense periodic orbits, the axiomatic predications and postulation alcovishness of chaos theory. Like gravity is a characteristic of space, so is tamas, rajas, sattva, with spaces knotted together by quantum information (for models see below) .accidents are preplanned actions of Terrestial and celestial Brahman AntiBrahman. All actions are perpetrated by them, be it train on fire or death. Belief systems play an important role in the determination of tamas, rajas, and satva state of human beings. There is neither cause nor effect. All actions are performed by Terrestial and celestial Brahman anti Brahman and it is off the record. psychic energy is negative equivalent to –mc^2, so that second law of thermodynamics is satisfied on a holistic basis vis-à-vis material energy

ParaBrahaman is infinite individual, collective and cosmic consciousnesses are infinite (concomitant general ledgers) for they are in the evolution mode with respect to each other in that order. Later emanates from the former as transactions emanate from the conservative state of assets and liabilities. Is not zero true for infinite conservative systems inside space and time? zero is the only consummation, consolidation and conservation. if fifty out of seventy members come out of the room at whatever time you come out, talk about the grand limbo of oblivion and hibernation of lifeless world you are going to and if you are not a celebrity, then it is not coincidence, it is coordination. Nature is playing the same trick on you. That a generalised law is there and you can measure does not make things less good albeit recalcitrant, refractory and rogue particles occur time and again, by and large. Entering in to all living beings through his constant part on the new moon night, he is born therefrom next morning. Absolute ParaBrahaman (zero) projects himself again and again on the screen of individual consciousness which appears of depairing despondent world. Some where the balancing act is done on karma. cosmic general ledger reflects upon itself to keep all transactional ties conservative (Brihadarnkayaka Upanishad: 1.5 14) for models see next paper through its divine power ( i interpret man creates god; it is attribution and ascription of various Shaktis would the formation of god takes place like space and time) the self assumes (eb) the form of a deity man can contemplate and venerate, even though siva, the pure subject, can never in fact be an object of meditation (adhyeya). Until we realise our true identity with Sankara, he is worshipped and conceived to be a reality alien to ourselves. While we are in the realm of creation, he too is a creation or mode or appearing of the absolute, manifest to us in meditation, through his freedom as an eternal, omniscient being. There is no gulf between the created and the uncreated creator: nothing in reality, although an object of knowledge, ceases to be Siva: this is the reason why meditation [on this or that aspect] of reality bestows its fruit. The world of the senses and mind appears to the well awakened (suprabuddha) as a theophany, an eternal revealing of god in his creation. The doctrine of vibration declares that "there is no state in word, meaning or thought, either at the beginning, middle or end, which is not Siva." 59 to utter any word is, in

reality, to intone a sacred formula.every act is a part of Siva's eternal cosmic liturgy, every movement of the body a ritual gesture (mudras), and every thought, god's thought. By what path are you not attainable? What words do not speak of you? in which meditation are you not an object of contemplation? What indeed are you not, o lord? 61 spandashakti, which accounts for the appearing of all things, is also the means by which deity in its many varied forms appears to man. Ksemaraja concludes: the ultimate object of worship of any theistic school differs not from the Spanda principle. The diversity of meditation is due solely to the absolute freedom of Spanda. 62 Sankara is not only the supreme object of devotion; as the static polarity of the absolute, he is the inner reality which holds together its Siva and Sakti outer manifestations. Phenomena are patterns of cognitions projected onto the surface of self-luminous siva-consciousness. There they become apparent, directly revealed to consciousness according to their manifest form. Siva is accordingly symbolized as the ground or surface of awareness, smooth and even like a screen (samabhittitalopama). 64 inscribed on this screen (kutfya) are the countless manifest forms which appear within it rendering it as diverse and beautiful as a fossil ammonite (falagrama). 65 Siva is the sacred ground upon which the cosmic mandalas are drawn, the absolute surface of inscription which bears the mark (cihna) of the universe. Abhinava writes: the variety of this world can only be manifest if the highest lord, who is essentially the pure light of consciousness, exists; just as a surface is necessary for a picture. if external objects were perceived in isolation then, because 'blue' and 'yellow', etc., are self-confined and the perceptions [we have of them] refer to their objects alone and so are insentient, mute and dumb in relation to one another . . . how would it be possible to be aware that an object is variegated? but just as depths and elevations can be represented by lines on a smooth wall, and we perceive [a female figure and think], similarly it is possible to be aware of differences in the variegated (contents of experience) only if all the diverse perceptions are connected together on the one wall of the universal light of consciousness. 66 Siva is the perfect artist who, without need of canvas or brush, paints the world pictures. The instant he imagines it, it appears spontaneously, perfect in every respect. The colours he uses are the varying shades

and gradations of his own Spanda energy and the medium his own consciousness. Intractable disaffect, Mutinous obdurate, Rebellious recalcitrant, Seditious refractory, Insubordinate insurgent, antagonistically antipodean, discrepant incompatible, Contumacious coadjutant, Conflicting converse, Unyielding obdurate, the antagonistic and anachronistic dispositional ties of Anti Brahman are in fact the sine qua non of Brahman only (See Vishwaroopa of Mahavishnu, which contains the polarised positives on RHS and negatives on LHS). The universe is coloured with the dye of its own nature (svabhava) by the power of Siva's consciousness (a//). Rajanaka rama says: homage to him who paints the picture of the three worlds, thereby displaying in full evidence his amazing genius {pratibha)\ to Sambhu who is beautiful with the hundreds of appearances laid out by the brush of his own unique, subtle and pure energy. 68 analogously, at the microcosmic level, all the cognitions and emotions, etc., which make up the individual personality form the outward flow of essentially introverted consciousness. They are specific pulsations (vise$aspanda) or aspects of the universal pulsation (samanyaspanda) of pure t consciousness. at the lower level, within the domain of maya, they represent the play in the fettered soul of the three primary qualities (gunas) or feeling-tones' which permeate to varying degrees his daily experience. These are: 1) sattva — the quality of goodness and luminosity which accompanies blissful experience both aesthetic and spiritual. 2) Rajas — the passion or agitation which oscillates between the extremes of 'light' and 'darkness' and characterises inherently painful experiences. 3) Tamas — the torpor and delusion which accompany states of inertia and ignorance. 69 the liberated soul recognises that these three are the natural and uncreated powers of pure consciousness. for him they are manifest respectively as: 1) Sankara 's power of knowledge (jnana) — the light of consciousness (prakasa); 2) the power of action (kriya) — the reflective awareness of consciousness (vimarsa)\ 3) the power of maya — which does not mean here the world of diversity, but the initial subtle distinction which appears between subject and object in pure consciousness. 70(doctrine of vibration: mark S.G. Dyczkowski) .to "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new

concepts and meanings that could not be included in the old system. at the heart of western metaphysics, for example, derrida finds the opposition between "speech" and "writing." this binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior in derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. The postmodern turn in philosophy: theoretical provocations and normative deficits by Steven best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html. You cannot see beyond consciousness for that is the limit of your storage of information or the basis for understanding further information. Turiya is empty bucket; it is from which prajna the measure of  arises; turiya measures the number of quantum information stores in other three states of waking, dreams, and dreamless deep sleep (Manduka Upanishad; my interpretation for models next paper).let there be path where angels tread: it is a force never born never dies the brahman an anti brahman verily there are parallel universes in (eb) in the ambit of space and time; the society of sinners and the sinned with the grammar of evil; parallel universes contain(e) multiverses in them; all physics, philosophy,  break down(e) in these universes; astronomical gratification ,deprivation, glorification, mortification, projections, axes, dimensions, rotations, foldings, concurrences, motions, shocks, are the axiomatic predications of multiverses and parallel universes; death is when self (the witness consciousness; register of individual general ledger) leaves the body; ostensibly parallel seem to have (e) an objective; parallel universes donot (e) have objectives; their paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common

attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, affection- appreciation, awareness- discernment, emotion-feeling, gut reaction, heart insight, intuition, judgment, keenness, perceptiveness, rationale, sensation, sense, sensitiveness, sensitivity, sentiment, susceptibility, taste, vibes, admiration, aesthetic sense, affection, appraisal, assessment, attraction, awareness, cognizance, commendation, comprehension, enjoyment, esteem, estimation, grasp, high regard, knowledge, liking, love, perceptional realization, recognitional regard, relish, respect, responsiveness, sensibility, sensitiveness, sensitivity, sympathy, understanding, valuation, acknowledgment, airing, baring, , confessional defenselessness, denudation- denunciation, disclosure, display, divulgence-exhibition, hazard, introduction, jeopardy, laying open, liability, manifestation, nakedness, openness, peril, presentation, publicity, revelation, risk, showing, susceptibility , susceptiveness,susceptivity, unfolding, unmasking, unveiling, vulnerability, vulnerableness, action, affection, appreciation and ardor, behavior, capacity, compassion, concern, cultivation and culture, delicacy, discernment and discrimination, emotion, empathy, faculty, fervor, fondness, heat, imagination, impression and intelligence, intensity and intuition, judgment, keenness- palpability, passion, pathos, pity, reaction, refinement, sensibility, sensitivity, sentiment, sentimentality , sharpness, spirit, sympathy, tangibility, taste, tenderness, understanding, warmth, fond memories, hearts and flowers, homesickness, longing, pining, reminiscence, ,schmaltz, sentimentality wistfulness, yearning satiation is only ostensible; their job is to get the job done; their business is to mind other's business ;all laws of physics and finance breakdown; it is antilaw,piscatorial Piratish and bubonic bucaneerishness; there is free transactions between in multiverses; when reflected on the screen of consciousness they appear as extant existential universe. Transactions are astronomical and gratification

deprivation, glorification and mortification incalculable; they have no barriers of region, religion, language which they create for others. theirs objectives are: corruption criminality, debasement debauchery, degradation evil immoral lewdness licentiousness, perversion sinfulness, vice viciousness, vileness wickedness, improbity corruption, atrocity decadence, degeneration degradation, depravity evil, immorality impurity, infamy iniquity, looseness lubricity, perversion profligacy, sinfulness turpitude,vice viciousness, vulgarity wickedness,disgrace bad reputation, abasement abuse, baseness black eye, blemish blur, brand comedown,contempt contumely, corruption culpability, debasement debasing, defamation degradation, derision disbarment, discredit disesteem, disfavor dishonor, disrepute disrespect, humbling humiliationignominy ill repute, infamy ingloriousness, meanness obloquy, odium opprobrium, pollution ,put-down reproach, scandal scorn, slander slight, slur spot, stain stigma, taint tarnish, turpitude venality, enormity Horribleness , abomination atrociousness, atrocity crime and the positive aspectionalities thereof .let there not be any incursions where evil mercenaries in the garb of mendicants thrive with adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatious gainsaying or conformingly pugnacious Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, execrable and goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness . statement for

Mephistophelean mercenaries philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality. Negative energy stabilising wormholes are considered candidates for (er) darkmatter. An analysis of negative energy might be a useful ingredient of the thoughts of travel at speeds more than the velocity of light. Peirce explains that an "imperfect continuum" possesses "topical singularities, or places of lower dimensionality" which break continuity, while "the parts of a perfect continuum have the same dimensionality

as the whole." (CP 4.642) Circular continua without singularities are time and space. (CP 6.210-212) By using math as his lens, Peirce envisioned singularities or lower dimensional holes in the imperfect cosmologic continuum, while maintaining that a probabilistic law of chance exists that entropy decreases **(Wikipedia).** On the side of philosophy**,** quite explicitly, Deleuze provides a critique of Husserl's "sleight-of-hand" change from a static categorical phenomenology to one supposedly more dynamically constitutive of reality. This a clear cut **chagrined circumvention**, **contraventional curbing**, **defeated disgruntlement**, **dissatisfactional downer**, **dragged failure**, **fizzled foil**, **grievance hindrance**, **impedimental irritation**, letdown, nonfulfillment, nonsuccess, obstruction, old one-two, **resent mental setback**, unfulfillment of the earlier theories makes Sartre contend that human existence is a conundrum whereby each of us exists, for as long as we live, within an overall condition of nothingness (nothing-ness)—that ultimately allows for free consciousness. But simultaneously, within our being (in the physical world), we are constrained to make continuous, conscious choices. It is this dichotomy that causes anguish, because choice (subjectivity) represents a limit on freedom within an otherwise unbridled range of thoughts. Subsequently, humans seek to flee our anguish through action-oriented constructs such as escapes, visualizations, or visions (such as dreams) designed to lead us toward some meaningful end, such as necessity, destiny, determinism (God), etc. Thus, in living our lives, we often become unconscious actors—Bourgeois, Feminist, Worker, Party Member, Frenchman, Canadian or American—each doing as we must to fulfill our chosen characters' destinies. However, Sartre contends our conscious choices (leading to often unconscious actions) run counter to our intellectual freedom. Yet we are bound to the conditioned and physical world—in which some form of action is always required. This leads to failed dreams of completion, as Sartre described them, because inevitably we are unable to bridge the void between the purity and spontaneity of thought and all-too constraining action; between the being and the nothingness that inherently coincide in our self. Deleuze insists on the absolute non-resemblance between what conditions (metastable, pre-individual plane of singularities) and what is conditioned

(actualized or the individuated 'thing'). Only then do "the conditions of the true genesis become apparent," Deleuze writes (Logic of Sense 105). Without addressing this issue, Deleuze's borrowings from Simondon, Lautman and structuralism remain untethered. Because of Bowden's decision to completely sidestep Deleuze's overt negotiations with phenomenology, his otherwise sophisticated commentary suggests a pedagogical narrowness. If, the omission of phenomenology strengthens Bowden's focus on the ontological priority of the event over substantial ontologies via the genealogy of figures and movements Deleuze outlines, it also incompletely articulates the problem of the sense-event as Deleuze formulates it. We have lost any idea of whom Deleuze takes to be his primary adversary, against whom he stakes out his position. This worry is most clearly manifested in Bowden's description of Deleuze's ontological prioritizing of events as offering a "transcendental ontology." For Bowden Deleuze's transcendental ontology defines events that "are ontologically prior to substances, 'all the way down'" (82). But Bowden provides a second definition: it is that the priority of events over worldly individuals means that the world of events "is not something external to the conditions of knowledge" (69). **SEAN BOWDEN the Priority of Events: Deleuze's Logic of Sense Sean Bowden, the Priority of Events: Deleuze's Logic of Sense, Edinburgh University Press, 2011, 296pp., $40.00 (hbk), ISBN 9780748643646. Reviewed by David Scott, Coppin State University Weltschmerz, agony, apprehension, blues, depression, dread, mid-life crisis, misgiving, nervousness, uneasiness, angst were expressed with such effulgent words…… aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement, argument, articulational assertion, asseverational choice of words, commentary and communication, declaration and definition, delivery, diction, elucidation, emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation about the human predicament is brought out in Being and Nothingness. Being and Nothingness offers a critique of Sigmund Freud's theory of the unconscious, based on the claim that consciousness is essentially self-conscious. Sartre also argues that**

**Freud's theory of repression is internally flawed. According to Sartre, in his clinical work, Freud encountered patients who seemed to embody a particular kind of paradox—they appeared to both know and not know the same thing. In response, Freud postulated the existence of the unconscious, which contains the "truth" of the traumas underlying the patients' behavior. This "truth" is actively repressed, which is made evident by the patients' resistance to its revelation during analysis. Yet what does the resisting if the patients are unaware of what they are repressing? Sartre finds the answer in what Freud calls the "censor". "The only level on which we can locate the refusal of the subject," Sartre writes, "is that of the censor.** The Holographic Principle, as conceived in current physics, applies to fields, and perhaps, even to more elementary entities, called strings and branes. These more elementary entities remain hypothetical at the time of this writing. There are many layers between the level of fields and that of neurons. We know that neurons are surrounded by a surface phospholipid membrane, which supports an electrochemical process that is fundamentally necessary to human experience. As the information supported by the membrane surface according to the Holographic Principle is proportional to its area, we should expect to find that the neural correlates of consciousness most clearly associated with those parts of the neuron that have the have a high proportion of neuronal surface area, and, in particular, receptive surface area. Dendrites account for an average of about 90% of the neuron is receptive surface area (Wong, 2002), and so would be expected to be the most important neural structures in information processes. We find this to be true **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science (Wikipedia).** Discordant-discrepant, dissonant-hard-line, implacable, incompatible, incongruous- inconsistent, inexorable-inflexible, inharmonious-intransigent, opposed, reluctant, unappeasable, uncompromising, unfriendly some of these expatiations and enumerations might sound, a more generalised interpretations is given by a Vedantin. Look at this arresting passage: The teacher continued to be silent. When addressed a second and third time he said: "I am teaching, but you do not follow. The Self is silence." The undetermined and unthinkable character of the Brahman is

a consequence of the absolute's eternal and immutable nature. To concede the existence of a real universe is, from the Vedantin's point of view, to posit the existence of a reality apart from the Brahman. Nor can we simply identify a real universe with the absolute unless we are prepared to compromise its unchanging, absolute status. The criterion of authenticity is immutability. Reality never changes; only that which is less than real can appear to do so. Reality is constant in the midst of change. What this means essentially is that there is change although nothing changes. This impossible situation is reflected in the ultimate impossibility of change itself. That which does not exist prior to its changing and at the end, after it has changed, must be equally non-existent between these two moments. Although the world of change appears to be real, it cannot be so. Change, according to the Vedantin, presupposes a loss of identity. Reality cannot suffer transformation; if it were to do so, it would become something else and the real would be deprived of its reality. The immortal can never become mortal, nor can the mortal become immortal. The ultimate nature of anything cannot change. Change of any sort is merely apparent (vivarta); the world of change and becoming is a false super- imposition (adhyaropa, adhyasa) on the absolute. In cosmic terms, the mistake (bhranti) consists of the supposition that the real Brahman is the unreal universe and the unreal universe is the real Brahman. In microcosmic terms, it is the mistake of falsely conceiving the body, mind or even one's personality to be the Self. In the same way as the image of a snake is falsely superimposed on a rope, similarly the universe is falsely projected onto the real substratum, the Brahman. Ignorance is not merely a personal lack of knowledge, but a cosmic principle. As such it is called "Maya," the indefinable factor (anirvacaniya) that brings this mistake in identity about. The reality status of this cosmic illusion is also indefinable: on the one hand it is not Brahman, the sole reality; on the other hand it is not absolutely non-existent like a hare's horn or the son of a barren woman. Brahman is the source of world appearances only in the sense of being their unconditioned ground or essential nature. The universe is false not because it has no nature of its own but because it does have one. Just as the illusion of a snake disappears when one sees that it is nothing but a rope, similarly cancellation (badha) of the empirically real occurs

when the absolute reality of the Brahman is realised. Thus, according to Vedanta, appearance implies the real, while the real need not imply appearance. To appear is essentially to appear in place of the real, but to be real is not necessarily to appear. All things -exist because the absolute exists. It is their Being. Thus the very existence of phenomena implies their non-existence as independent realities.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^{\infty} \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9 \left( U_9 - V_9 \int_0^{\infty} cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9 \right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9 \int_0^{\infty} cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$then given the constraints of $f(T_9), (A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 -$$

$$V_9 M_9 -]e^{-T_9\lambda_9} = 0$$

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9 Z_9 \sin(Z_9 T_9) - (A_9 M_9 + Y_9) + (A_9 M_9 + X_9)\cos(Z_9 T_9) = 0$$

$$(A_9 M_9) Z_9 - M_9 Z_9 \cos(Z_9 T_9) + (A_9 M_9 - X_9)\sin(Z_9 T_9) = 0$$

Quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High

Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) **Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP** This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) **Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP** In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference: Journal of Physics: Conference Series 70 **Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows:** System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol

Where $A_9 = U_9 - V_9$Using (A2.3), substitute for $\boldsymbol{Z}_9^2$ and $Z_9$in (A2.2):

$W_9^2 (X_9^2 - 3 M_9 A_9 X_9 - M_9^2 (X_9 - Y_9) + 2 M_9^3 A_9 + 2 A_9^2 M_9^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_9$=1-cos ($Z_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (X_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ X_9 - Y_9/2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \leq W_9 \leq 2$ *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (X_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\,M_9)^2\,(2A_9 M_9 - (X_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9 M_9^2 - 3M_9 X_9\right) + W_9\left(3M_9 A_9^2 - 2A_9 X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9+Y_9\right)$$

$$=2M_9^2W_9^2(M_9+\mathrm{A}_9)+M_9A_9^2W_9+\left(2A_9M_9-X_9+Y_9\right)(A_9W_9+M_9W_9^2$$

$$+(2-W_9)((2M_9X_9+M_9Y_9)W_9+A_9(X_9+Y_9)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range 0 $W_9$ 2.Defining

$Q_9\left(W_9,A_9\right)=a_9W_9^2+c_9$ *where* $a_9b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

To Heidegger writings we shall compose the model. Following is not still written model able form. For one of the models please see any one of the papers under the appellation: Ïn Excited Reveries

In his later writings Heidegger drops the term 'phenomenon' in preference for the verbal form 'phainesthai' in order to emphasize even more the actuality or presentational property of Being. Explaining this new form of the term he writes: "Being disclosed itself to the ancient Greeks as 'physis'. The etymological roots 'phy-' and 'pha-' designate the same thing: 'phyein', the rising-up or upsurge which resides within itself as 'phainesthai', lighting-up, self-showing, coming-out, appearing-forth." Heidegger contrasted his notion of phenomenon with semblance (Schein) and with appearing (Erscheinung). In the case of semblance a thing can show itself as that which it is not, as when fool's gold shows itself to be gold. The ancients always allied semblance with non-being. Heidegger points out, however, that semblances are grounded in

showings, and so does Abhinava. Both Heidegger and Abhinava consequently maintain that all semblances have a real basis and are to be treated as instances of phenomena along with the so-called real showing or manifestation of non-deceptive objects. So Heidegger states that: 'how- ever much seeming, just that much being'. Thus self-showing or appearing defines Being as phenomenon, but this definition of Being is as yet incomplete. Being is not only self-showing but 'logos' which Heidegger explains means 'discourse' (Rede) in the sense of 'apophansis': 'letting-be-seen'. Phenomenology, which according to Heidegger is the only correct study of Being, means 'letting-be-seen-that-which- shows-itself. This is true of Saiva Paramadvaita as well. The reality of the world demands recognition; we are forced to accept the direct presentation of the fact of our daily experience. As Abhinava says: "if practical life, which is useful to all persons at all times, places and conditions were not real, then there would be nothing left which could be said to be real" A thousand proofs could not make 'blue' other than the colour blue. The reality of whatever appears in consciousness cannot be denied. Objects appear; they do not cease to do so by a mere emphatic denial. The manifestation of an entity in its own specific form is a fact at one level of consciousness; it is real. The appearing of the same entity in the same form but recognised to be a direct representation of the absolute is also a fact, but at another level of consciousness. It is no more or less real than the first. 'As is the state of consciousness, so is the experience,' says Abhinava. Although the nature of the absolute is discovered at a higher level of consciousness, The Doctrine of Vibration nonetheless it presents itself to us directly in the specific form in which we perceive things; otherwise there would be no way in which we could penetrate from the level of appearing to that of its source and basis. Abhinava writes: Real is the entity iyastu) that appears in the moment of direct perception (sak$atkara), that is to say, within our experience of it. Once its own specific form has been clearly determined one should, with effort, induce it to penetrate into its pure conscious nature. All things are known to be just as they present themselves. The concrete actuality of

being known (pramiti), irrespective of content, is itself the vibrant (spanda) actuality of the absolute. Liberating knowledge is gained not by going beyond appearances but by attending closely to them. "The secret," Mahesvarananda says, "is that liberation while alive (jivanmukti) is the profound contemplation of Maya's nature." No ontological distinction can be drawn between the absolute and its manifestations because both are an appearing (dbhasa), the latter of diversity and the former of 'the true light of consciousness which is beyond Maya and is the category Siva'. Those who have attained the category of Pure Knowledge above Maya and have thus gone beyond the category of Maya, see the entire universe as the light of consciousness . . . Just as the markings [on a feather] are nothing apart from the feather, the feather [is nothing apart from] them, similarly, when the light of consciousness is manifest, the whole group of phenomena is manifest as the light of consciousness itself. Within the sphere of Maya, every entity's 'own nature' (svabhdva) corresponds to its specific manifest form. Accordingly it is defined as that which distinguishes it from all else and from which it never deviates. Above the sphere of Maya, that is, above the level of objectivity, is the domain of the subject. At this level, everything is realised to be part of the fullness of the experience r and hence no longer bound by the conditions which impinge on the object. Here the part is discovered to be the whole, that is, consciousness in toto. In this sphere beyond relative distinctions, the yogi realises that (all) the categories of existence are present in every single category. m The yogi experiences every individual particular as the sum total of everything else. He recognises that all things have one nature and that every particular is all things. This is the 'essence' (sard) or co-extensive unity (samarasya) of all things. Integral Monism of Kashmiri SaivismWe have established that reality is manifest according to how [and the degree in which] the freedom of consciousness reveals it and that [this freedom] is the womb of all forms. Just as 'sweetness' is present in its entirety in every atom of the sugarcane, so each and every atom [of the universe] bears within itself the emanation of all things. This is

the level of consciousness in which the absolute reflects on itself realising to its eternal delight and astonishment (camatkara) its own integral nature. The reality of the world of diversity is not denied, but experienced in a new mode of awareness free of time and space in the eternal omnipresence of the Here and Now. [Phenomenal forms of awareness] such as 'this [exists]', born of the colouring [imparted to the absolute] by the limitations engendered by the diversifying power of time {kalakalana) also emanate within the Supreme Principle. There [at that level], Fullness {purriata) is the one nature [of all things] and so everything is omnipresent; otherwise, associated with division (khancjlana), the Fullness [of the absolute] would not be full. The content of absolute consciousness consists of diverse appearings (abhdsa) which, because they are manifest through it in this way, do not compromise the wholeness of consciousness. Everything we perceive is a momentary collocation of a number of such manifestations which combine together like 4 a row of altar lamps' (dipavalT) to form the single radiant picture of the universe. The individual objects which constitute the universe are specific collocations of such 'atomic' appearings. Together they form a single unified particular which appears according to its own defining features (svalaksand). A jar, for example, consists of a number of appearances such as 'round', 4 fat\ 'earthen', 'red', etc., which together discharge a single function (arthakriya), in this case, that of carrying the appearance 'water'. They unite with each other much as the scattered rays of a lamp come together when focused, or as the various currents of the sea together give rise to waves. Atomic appearings can combine in any number of ways, provided that they are not contrary to one another as established by the dictates of natural law (niyati). An appearance of 'form', for example, cannot combine with that of 'air'. Insofar as they share a common basis (samanyadhikaranya), a given cluster of appearances appears as a single whole. This common basis is the most prominent member of the group; the appearance 'jar' is such in the example quoted above. Any one appearance in a cluster may assume a more important or subordinate role.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_9 \square - b_9 \square 0$$

And

$$b_9^2 \square 4a_9c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9 in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\ X_9$ ,

$Let\ Y_9 = \ y_9\ X_9$And try the substitution $A_9 = X_9/(2 + A_9\ y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + \ n_9 y_9)$ we have

$$W_9^2\ y_9\left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^2\left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9\,X_9\left[\frac{X_9\left(1+n_9y_9\right)}{M_9^2\left(2+n_9y_9\right)^2}+2y_9\left(n_9-1\right)+\frac{X_9^2\left(9+y_9\right)}{M_9^2\left(2+n_9y_9\right)}\right]=0.\ (A2.12)$$

Since $n_9$ and $A_9$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9=\left(2+n_9y_9\right),$ *the condition*$(A2.99?)$*can be written as*

$$\left[\frac{X_9\left(E_9-9\right)}{E_9^2}\right]^2\left[\frac{X_9^2}{M_9^2}-4y_9\,E_9^2\frac{X_9}{M_9^2}+4y_9\left(n_9-2\right)E_9^3\right]$$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9/W_9^2)$ can be expressed as

$$2E_9\left[y_9E_9+\left(y_9\,E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}\right]>\frac{X_9}{M_9^2}>2E_9\left[y_9E_9-y_9E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}]$$

dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique

for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) **Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP** Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) **Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP** Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a

cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) **Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP** Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service

placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) **Dynamic Service Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP** Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by $H*(X,(\wedge*TX)\otimes g)$, where X is the target hyperk\"ahler manifold. We further give a mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) **A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van**

**Raamsdonk1 and Nathan Seiberg2** study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. **Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP** Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy **The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP** Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the

present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. **ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP** Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition. Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb) the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537

© 2016 American Physical Society **Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP** Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant ($\Delta\alpha/\alpha$) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and $\Delta\alpha/\alpha$ to (e&eb) this conclusion. However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts $z\lesssim 3$, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the $\Lambda$ cold dark matter model value ($j\Lambda CDM=1$); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in $\alpha(z)$, linked to (e&eb) a linear coupling ($\zeta$) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order $|\zeta|\sim 10-4$, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in $\Delta\alpha/\alpha$ (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society **The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP** Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation

rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void

Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems, authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society **Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP** Authors measure the star formation properties of (e) two large samples of

galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the **star formation** rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a **function of environment**. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]**.** Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size $\sqrt{n} \times \sqrt{n}$? The problem is that $\sqrt{n}$ time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time $O(\sqrt{n})$ for $d \geq 3$, or $O(\sqrt{n}\log 5/2n)$ for $d = 2$. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search

tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ∗This work was mostly done while the author was a **PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship and by CIAR**. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 **Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP** Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can

lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 **Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP** Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. **Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP** Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e)

sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) **Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP** The relationship between the foundations of mathematics and (e&eb) physics is a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb) flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling

field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) **Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP** The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar

and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph] (or arXiv: 1503.05600v1 [math-ph] for this version) **Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP** This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs

mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) **Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP** Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and

derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to **the fine structure constant**. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) **Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP** Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor goes to 0 as the (e&eb) time goes to 0, the

big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph] (or arXiv: 1306.4613v1 [quant-ph] for this version) **Effects of gauge theory based number scaling on geometry Paul Benioff STOP** In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x,

outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph]

(or arXiv: 1211.3381v2 [quant-ph] for this version) **Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP** Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value,

\int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances, the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) **Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP** The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include

(e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) **Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP** Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) **New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP** This work is based on (e) the field of reference frames based on (e) quantum

representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects:Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc) Journal reference:Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) **A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul B**

**ABSOLUTE STABILITY**

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_9$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9-9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2-2y_9\frac{(E_9+n_9^2y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right)-4y_9(E_9-n_9)$$ (A2.14)

And has the solution

$$\frac{X_9}{M_9^2}>\frac{E_9}{(E_9-9)}[y_9(E_9+n_9^2y_9)+[y_9^2(E_9+n_9^2y_9)^2+4y_9(E_9-9)(E_9-n_9)]^{9/2}]$$

Given $E_9>n_9$this lower limit to $X_9/M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9/M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9<n_9$ $(ie.,2A_9M_9^2<X_9-Y_9)$then (A2.9)and (A2.8)show that $Q_9$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2}>2E_9[y_9E_9-(y_9E_9(9+y_9)(E_9-n_9)]^{9/2}]$$

And

$$E_9>n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2/2M_9(X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$ and $X_9$ and $Y_9$ by $W_9^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

**Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability

of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e) long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007 STOP. David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer P** examine the possibility of coherent reversible information transfer between solid-state superconducting qubits and (e&eb) ensembles of ultracold atoms. Strong coupling between these systems is mediated by (e&eb) a microwave transmission line resonator that interacts near resonantly with (e&eb) the atoms via (e&eb) their optically excited Rydberg states. Solid-state qubits can then be used to (e) implement rapid quantum logic gates, while (e) collective metastable states of (e) the atoms can be employed for (e) long-term storage and (e&eb) optical readout of quantum information. **Reversible state transfer between superconducting qubits and (e&eb) atomic ensembles David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer Phys. Rev. A 79, 040304(R) – Published 14 April 2009 STOP.** Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures have been proposed. An

attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. **Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier** demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in (eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. **Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP.** In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown

to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , **Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally** frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. **Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg** propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states

into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. **Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011 http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan** proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 **Qubit transfer between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller** propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. **Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys.**

**Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich** report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. **Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi** investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. **Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev** investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer.

Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. **L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP.** In recent years **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P have** investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P** apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys. Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. **Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S.**

**Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP.** Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert **Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe** realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.**220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2** [quant-ph] for this version) **Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP.** Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. **Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker** describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue

from over (e&eb) a thousand plant and animal samples. **NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP.** Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this behavior, enhancing and accelerating (eb, eb+) the growth of correlations. **Michael Kastner** gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. **Michael Kastner** tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) **Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP.** A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. **Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E. Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91,**

**075435 – Published 26 February 2015 STOP.** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still

problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to

(e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP** Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that

(eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035

**Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime**

**C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013**

**STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and

cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP** Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex

Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) **A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen**

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9 and$ separate (A3.1)into real and imaginary parts

$$Z_9^2 + \left(M_9 + V_9\right) Z_9 \sin\left(Z_9 T_9\right) - \left(U_9 M_9 + Y_9\right) + (M_9\left(U_9 + V_9\right) - X_9)$$

$$\mathrm{X}\cos\left(Z_9 T_9\right) - 2M_9 V_9 cos^2\left(Z_9 T_9\right) + M_9 V_9 = 0$$

$$(M_9 + U_9)\, Z_9 - (M_9 + V_9)\, Z_9\ cos\left(Z_9 T_9\right) + \left(M_9\left(U_9 + V_9\right) - X\right)$$

$$\mathrm{X}\sin\left(\left(Z_9 T_9\right) - M_9 V_9 \sin\left(Z_9 T_9\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 =$ $1 - cos\, Z_9\ T_9$and $A_9 = U_9 - V_9$we obtain

$$4\, M_9^2 V_9^2 W_9^4 + 2\, M_9 V_9\ W_9^3[(M_9 - V_9)^2 + A_9\left(M_9 + V_9\right) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2\left[M_9 - V_9\right)^2\left(2\, A_9\ M_9 + Y_9 - X_9\right) + \left(2A_9\ M_9 - X_9\right)^2 - A_9\left(M_9 - V_9\right)\left(2A_9\ M_9 - X_9\right)$$

$$+ M_9 V_9\,(A_9^2 + 4\left(X_9 + Y_9\right))]$$

$$+ W_9\left[A_9^2 + 2X_9\right)\left(2\, A_9\ M_9 - X_9\right) + 2A_9 Y_9\,(M_9 + V_9) - M_9 A_9^3] + A_9^2\left(X_9\ M_9 - Y_9\right) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9(2, U_9 V_9) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting cos $(Z_9 T_9) = -1\ in\ (A3.3)$

We have applied the model to Kashmiri Shaivism. Whose elaborate rituals is enough to make a indivual raise from individual consciousness to cosmic consciousness which we have called Cosmic Consciousness.

It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) **Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP** The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic

Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** The group of unimodular transformations on the roots of polynomials over GF(2) is considered, and those polynomials with (e&eb) symmetries in (eb) the unimodular group are identified. The cross-correlation function between (e&eb) two maximum-length linear shift register sequences of (e) the same degree is shown to be (=) computable as an explicit linear transformation, in (eb) matrix form, on (e&eb) either one of the sequences, regarded as (=) a vector. The underlying vector space is (=) the "cyclotomic algebra," generated by (e) the cyclotomic cosets, or by what Gauss termed the **"periods" of the cyclotomic equation**. A variety of numerical examples are worked in detail. **Information Sciences Volume 1, Issue 1, December 1968, Pages 87-109 Theory of transformation groups of polynomials over GF(2) with applications to linear shift register sequences Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more link. Solomon W. Golomb** stop Authors study the P−V criticality and phase transition in (eb) the extended phase space of charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of (e) the system and its conjugate quantity is (=) the thermodynamic volume of the black hole. The black holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb) no P−V criticality and phase transition appear, while (e) for the black holes with a spherical horizon, even when the charge of the black hole is absent, t(eb) he P−V criticality and the small black hole/large black hole phase transition will appear, but it happens only in d=5 dimensions; when the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to have (e) the P−V criticality

and (e&eb)the small black hole/large black hole phase transition, there exists (ebb) an upper bound for the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the **Gauss-Bonnet coefficient** and Q is (=) the charge of the black hole. Authors calculate the critical exponents at the critical point and find that (eb) for all cases, (eb) they are the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09(2013)005 Cite as: arXiv:1306.6233 [gr-qc] (or arXiv:1306.6233v4 [gr-qc] for this version) **P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in IR but with (e&eb) improved UV behaviors. Here, authors study an IR modification which breaks ``softly'' (e) the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Authors obtain the black hole and cosmological solutions for **``arbitrary'' cosmological constant** that represent (eb) the analogs of the standard Schwarzschild-(A)dS solutions which can be asymptotically (A)dS as well as flat and authors discuss their thermodynamical properties. They also obtain solutions for (e) FRW metric with (e&eb) an arbitrary cosmological constant. Authors study its implication to (e&eb) the dark energy and find that it seems to be consistent with (e&eb) current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP.** The Einstein equations with (e&eb) a negative cosmological constant admit (e) black hole solutions which are asymptotic to (e) anti-de Sitter space. Like black holes in asymptotically flat space, these solutions have (e) thermodynamic properties including (e) a characteristic temperature and intrinsic entropy equal to (=) one quarter of the area of the event horizon in Planck units. There are however some important differences from (e) the asymptotically flat case. A black hole in anti-de Sitter space has (e) a minimum temperature which occurs when (e) its size is of the order of the characteristic radius of (e) the anti-de Sitter space. For larger black holes the red-shifted

temperature measured at (eb) infinity is greater. This means that such black holes have (e) positive specific heat and can be in stable equilibrium with (e&eb) thermal radiation at a fixed temperature. It also implies (eb) that the canonical ensemble exists for asymptotically anti-de Sitter space, unlike the case for asymptotically flat space. One can also consider the microcanonical ensemble. One can avoid (e) the problem that arises in asymptotically flat space of (e) having to put the system in (eb) a box with unphysical perfectly reflecting walls because (e) the gravitational potential of anti-de Sitter space acts as a box of finite volume. Communicated by A. Jaffe **Communications in Mathematical Physics December 1983, Volume 87, Issue 4, pp 577-588 Thermodynamics of black holes in anti-de Sitter space S. W. Hawking, Don N. Page STOP** In previous work authors introduced the concept of black-hole entropy, which they identified with (e&eb) the surface area of the black hole in question expressed in units of (e) the Planck length squared. Authors suggested that the appropriate generalization of the second law for a region containing a black hole is (=) that the black-hole entropy plus (+) the common entropy in the black-hole exterior never decreases. Here authors establish the validity of this law for the infall of an entropy-bearing system into (e&eb) a much larger and more massive generic stationary black hole. To do this authors determine (eb) a general lower bound for the increase in black-hole entropy, and an upper bound for the entropy of the system, while allowing for quantum effects at each stage. In passing authors show that the generalized second law is a statistical law which becomes (=) over-whelmingly probable in the limit of a macroscopic system. Authors also consider briefly more general situations. Finally, we give two simple examples of predictions made by the generalized second law for (e) black-hole formation processes. Received 17 September 1973 DOI:http://dx.doi.org/10.1103/PhysRevD.9.3292 **Generalized second law of thermodynamics in black-hole physics Jacob D. Bekenstein Phys. Rev. D 9, 3292 – Published 15 June 1974 STOP** The Lovelock gravity consists of (e) the dimensionally extended Euler densities. The geometry and horizon structure of (e) black hole solutions could be quite complicated in (eb) this gravity; however, authors find that some thermodynamic quantities

of the black holes like the mass, Hawking temperature and entropy have (e) simple forms expressed in terms of (e) horizon radius. The case with black hole horizon being a Ricci flat hypersurface is (=) particularly simple. In that case the black holes are (=) always thermodynamically stable with (e&eb) a positive heat capacity and their entropy still obeys (e&eb) the area formula, which is no longer valid for (e) black holes with positive or negative constant curvature horizon hypersurface. In addition, for black holes in (eb) the gravity theory of Ricci scalar plus a 2n-dimensional Euler density with (e&eb) a positive coefficient, thermodynamically stable small black holes always exist in (eb) D=2n+1 dimensions, which are absent in (eb) the case without (e) the Euler density term, while (e) the thermodynamic properties of the black hole solutions with (e&eb) the Euler density term are qualitatively similar to those of black holes without (e) the Euler density term as D>2n+1. Copyright © 2004 Published by Elsevier B.V. **Physics Letters B Volume 582, Issues 3–4, 4 March 2004, Pages 237–242 A note on thermodynamics of black holes in Lovelock gravity Rong-Gen CAI STOP** There are a number of similarities between black-hole physics and (e, eb, =) thermodynamics. Most striking is the similarity in the behaviors of black-hole area and (e&eb) of entropy: Both quantities tend to increase (eb+) irreversibly. In this paper we make this similarity the basis of (e) a thermodynamic approach to black-hole physics. After a brief review of the elements of the theory of information, authors discuss black-hole physics from the point of view of information theory. Authors show that it is natural to introduce (eb) the concept of black-hole entropy as (=) the measure of information about (e&eb) a black-hole interior which is (=) inaccessible to an exterior observer. Considerations of simplicity and consistency, and dimensional arguments indicate (eb) that the black-hole entropy is equal to (=) the ratio of the black-hole area to the square of the Planck length times a dimensionless constant of (e) order unity. A different approach making use of the specific properties of Kerr black holes and of concepts from information theory leads to (eb) the same conclusion, and suggests (eb) a definite value for the constant. The physical content of (e) the concept of black-hole entropy derives from (e) the following generalized version of the second law: When common entropy goes down a

black hole, (eb) the common entropy in (eb) the black-hole exterior plus the black-hole entropy never decreases. The validity of this version of the second law is supported by (e) an argument from information theory as well as by several examples. Received 2 November 1972 DOI:http://dx.doi.org/10.1103/PhysRevD.7.2333 **Black Holes and Entropy Jacob D. Bekenstein Phys. Rev. D 7, 2333 – Published 15 April 1973 STOP** The effects of power-law Maxwell field on the Van der Waals like phase transition of higher-dimensional dilaton black holes are probed in detail. It is shown that the Smarr relation gains corrections due to (e) the effects of (e) both the power-law Maxwell field and (e&eb) the dilaton field while (e) thermodynamic volume is exactly the same as (=) that of Einstein-Maxwell-dilaton black holes. Authors successfully derive the analytic solutions of (e) critical point and carry out (e&eb) some check to ensure that these critical quantities are (=) positive. It is shown that the constraint on the parameters turns out to be (=) 0<\alpha^2<1, which is tighter than that in (eb) the non-extended phase space. It is also shown (eb) that these critical quantities and (e&eb) the ratio P_cv_c/T_c are affected by (e&eb) the power-law Maxwell field. Moreover, critical exponents are found to coincide with those of other AdS black holes, showing (eb) the powerful influence of mean field theory. Comments: references added Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1601.05500 [gr-qc] (or arXiv:1601.05500v2 [gr-qc] for this version **Effects of power-law Maxwell field on the Van der Waals like phase transition of higher dimensional dilaton black holes Jie-Xiong Mo, Gu-Qiang Li, Xiao-Bao Xu STOP** A consistent and unified picture for (e) critical phenomena of charged AdS black holes in f(R) gravity is drawn in this paper. Firstly, we investigate the phase transition in canonical ensemble. Authors derive the explicit solutions corresponding to (e&eb) the divergence of CQ. The two solutions merge into (e&eb) one when (e) the condition Qc=−13R0−−−√ is satisfied. The curve of specific heat for Q<Qc has two divergent points and can be divided into (e&eb) three regions. Both the large radius region and the small radius region are (=) thermodynamically stable with (e&eb) positive specific heat while (e) the medium radius region is unstable with (e&eb) negative specific heat. However, when Q>Qc, the specific heat is

always positive, implying (eb) the black holes are locally stable and no phase transition will take place. Secondly, both the T−r+ curve and T−S curve f(R) AdS black holes are investigated and they exhibit Van der Vaals like behavior as the P−v curve in the former research. Critical physical quantities are obtained and they are consistent with (=) those derived from the specific heat analysis. Authors carry out numerical check of Maxwell equal area law for the cases Q=0.2Qc, 0.4Qc, 0.6Qc, 0.8Qc. The relative errors are (=) amazingly small and can be negligible. So the Maxwell equal area law holds for (e) T−S curve of f(R) black holes. Thirdly, authors establish geometrothermodynamics for (e) f(R) AdS black hole to examine the phase structure. It is shown that the Legendre invariant scalar curvature R would diverge exactly where (e) the specific heat diverges. To summarize, the above three perspectives are consistent with (e&eb) each other, thus providing (eb) a unified picture which deepens the understanding of (e) critical phenomena of (e) f(R) AdS black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1602.01251 [gr-qc] (or arXiv:1602.01251v1 [gr-qc] for this version) **A consistent and unified picture for critical phenomena of f(R) AdS black holes Jie-Xiong Mo, Gu-Qiang Li, Yu-Cheng Wu STOP** The state equation of the charged AdS black hole is reviewed in (e&eb) the T−r plane. Thinking of the phase transition, the T−S, P−V, P−v graphs are plotted and then the equal area law is used (e) in the three cases to get (eb) the phase transition point (P, T). The analytical phase transition point relations for (e) P-T of charged AdS black hole has been obtained successfully. By comparing the three results, authors find that the equal area law possibly cannot (e) be used directly for P−v plane. According to the T−S, P−V results, they plot (e&eb) the P−T−Q graph and find that for a highly charged black hole a very low temperature condition is required for (e) the phase transition. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI:10.1140/epjc/s10052-015-3641-0 Cite as: arXiv: 1503.07658 [gr-qc] (or arXiv:1503.07658v1 [gr-qc] for this version) **A note on Maxwell's equal area law for black hole phase transitionShan-Quan Lan, Jie-Xiong Mo, Wen-Biao Liu STOP** Recently, extended phase space thermodynamics of Lovelock AdS black holes has been of great interest. To provide

insight from (e) a different perspective and gain (e) a unified phase transition picture, non-extended phase space thermodynamics of (e) (n+1)-dimensional charged topological Lovelock AdS black holes is (=) investigated detailedly in (eb) the grand canonical ensemble. Specifically, the specific heat at (eb) constant electric potential is calculated and phase transition in the grand canonical ensemble is discussed. To probe the impact of the various parameters, they utilize (e) the control variate method and solve (e) the phase transition condition equation numerically for the case k=1, −1. There are two critical points for (e) the case n=6, k=1 while there is only one for other cases. For k=0, there exists no phase transition point. To figure out the nature of phase transition in the grand canonical ensemble, authors carry out an analytic check of the analog form of (e) Ehrenfest equations proposed by Banerjee et al. It is shown that Lovelock AdS black holes in (eb) the grand canonical ensemble undergo (e&eb) a second order phase transition. To examine the phase structure in (eb) the grand canonical ensemble, authors utilize the thermodynamic geometry method and calculate both the Weinhold metric and Ruppeiner metric. It is shown that for both analytic and graphical results that the divergence structure of (e) the Ruppeiner scalar curvature coincides with (=) that of the specific heat. Research provides (eb) one more example that Ruppeiner metric serves as (=) a wonderful tool to (e) probe the phase structures of (e) black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C75 (2015) 5, 211 DOI:10.1140/epjc/s10052-015-3434-5 Cite as: arXiv: 1503.01956 [gr-qc] (or arXiv: 1503.01956v1 [gr-qc] for this version **Non-extended phase space thermodynamics of Lovelock AdS black holes in grand canonical ensemble Jie-Xiong Mo, Wen-Biao Liu STOP** To provide an analytic verification of the nature of phase transition at the critical point of P−V criticality, the original expressions of Ehrenfest equations have been introduced directly. By treating the cosmological constant and (e&eb) its conjugate quantity as (=) thermodynamic pressure and volume respectively, authors carry out analytical check of (e) classical Ehrenfest equations. To show that approach is universal, authors investigate not only higher-dimensional charged AdS black holes (e&eb), but also rotating AdS black holes. Not only are the examples

of Einstein gravity shown, but also the example of modified gravity is presented for (e) Gauss-Bonnet AdS black holes. The specific heat at constant pressure CP, the volume expansion coefficient α and the isothermal compressibility coefficient κT are found to diverge exactly at the critical point. It has been verified that both Ehrenfest equations hold at (eb) the critical point of P−V criticality in (eb) the extended phase spaces of (e) AdS black holes. So the nature of the critical point of P−V criticality of AdS black holes has been demonstrated analytically to be (=) a second-order phase transition. These results are consistent with (e&eb) the nature of liquid-gas phase transition at the critical point. In this sense, extant research would deepen the understanding of (e) the relations of AdS black holes and (e&eb) liquid-gas systems. Moreover, successful approaches to introduce (eb) the original expressions of Erhenfest equations directly into (e&eb) black hole phase transition research demonstrate again (eb) that black hole thermodynamics is closely related (e&eb) to classical thermodynamics, which allows (eb) us to borrow techniques from (e) classical thermodynamics to investigate the thermodynamics of black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 89, 084057 (2014) DOI: 10.1103/PhysRevD.89.084057 Cite as:arXiv: 1404.3872 [gr-qc] (or arXiv:1404.3872v1 [gr-qc] for this version) **Ehrenfest scheme for P−V criticality of higher dimensional charged black holes, rotating black holes and Gauss-Bonnet AdS black holes Jie-Xiong Mo, Wen-Biao Liu STOP** Aiming at a unified phase transition picture of (e) the charged topological black hole in (eb) Ho\v{r}ava-Lifshitz gravity, authors investigate this issue not only in (eb) canonical ensemble with (e&eb) the fixed charge case but also in grand-canonical ensemble with (e&eb) the fixed potential case. Authors firstly perform the standard analysis of the specific heat,(e&eb) the free energy and the Gibbs potential, and (e&eb) then study its geometrothermodynamics. It is shown that the local phase transition points (eb) not only witness the divergence of (e) the specific heat, but also witness the minimum temperature and the maximum free energy or Gibbs potential. They also witness the divergence of (e) the corresponding thermodynamic scalar curvature. No matter which ensemble is chosen, the metric constructed can

successfully produce (eb) the behavior of the thermodynamic interaction and (e&eb) phase transition structure while other metrics failed to (e) predict the phase transition point of (e) the charged topological black hole in former literature. In grand-canonical ensemble, authors have discovered the phase transition which has not been reported before. It is similar to (e) the canonical ensemble in which the phase transition only takes place when k=−1. But it also has its unique characteristics that the location of the phase transition point depends on (eb) the value of potential, which is different from (e) the canonical ensemble where (e) the phase transition point is (=) independent of the parameters. After an analytical check of Ehrenfest scheme, authors find that the new phase transition is (=) a second order one. It is also found that the thermodynamics of (e) the black hole in Horava-Lifshitz gravity is quite different from that in (eb) Einstein gravity. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP1310(2013)056 DOI:10.1007/JHEP10(2013)056 Cite as: arXiv:1404.2497 [gr-qc] (or arXiv:1404.2497v1 [gr-qc] for this version) **A unified phase transition picture of the charged topological black hole in Horava-Lifshitz gravity Jie-Xiong Mo, Xiao-Xiong Zeng, Gu-Qiang Li, Xin Jiang, Wen-Biao Liu STOP** To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case $\beta\to\infty$ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness condition of (e) entropy. It is shown

that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure authors have discovered. It is shown that Legendre invariant thermodynamic scalar

curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws, leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions, geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged AdS black holes by considering the case of the RN-AdS black holes

surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase transition. For the case $\omega_q=-2/3$ authors derive analytic expressions of (e) critical physical quantities, while (e) for cases $\omega_q\neq-2/3$ appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this

version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=) chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy

Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed

quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions, while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma?**

All this consists of the act of consciousness alone", says Vasubandhu, "because unreal entities appear, just as a man with defective vision sees unreal hair or a moon, etc."He points to dreams as examples of purely subjective constructs which appear to be objective realities. The apparent reality dreams possess is not derived from any concrete, objective world, but merely from the idea of objectivity. While the Yogacara does not say that an idea has, for example, spatial attributes, it does have a form manifesting them. While he agrees with the Saiva idealist that

appearances have no independent existence apart from their appearing to consciousness, he maintains that for this reason they are unreal. The creativity of consciousness consists in its diversification in many modes having apparent externality; it is not a creation of objects. While the Kashmiri Saivite agrees that the world is pure conscious- ness alone, he maintains that it is such because it is a real creation of consciousness. The effect is essentially identical with the cause and shares in its reality. Matter and the entire universe are absolutely real, as 'congealed' (sty ana) or 'contracted' (samkucita) forms of consciousness. "This God of consciousness", writes Ksemaraja, "generates the universe and its form is a condensation of His own essence (rasa)" m By boiling sugarcane juice it condenses to form treacle, brown sugar and candy which retains its sweetness. Similarly, consciousness abides unchanged even though it assumes the concrete material form of the five gross elements. The same reality thus abides equally in gross and subtle forms. Consequently no object is totally insentient. Even stones bear a trace (vasana) of consciousness, although it is not clearly apparent because it is not associated with the vital breath (prana) and other components of a psycho-physical organism. Somananda goes so far as to affirm that physical objects, far from being insentient, can only exist insofar as they are aware of themselves as existing. The jar performs its function because it knows itself to be its agent. Indeed, all things are pervaded by consciousness and at one with it and hence share in its omniscience. Thus, Siva, Who perceives Himself in the form of physical objects, is the one ultimate reality. "The jar knows because it is of my nature", writes Somananda, "and I know it because I am of the jar's nature. I know because I am of Sadasiva's nature and He knows because He is of my nature; Yajnadatta [knows] because he is of Siva's nature and Siva [knows] because He is of Yajnadatta's nature". Integral Monism of Kashmiri Saivism Everything in this sense is directly perceived by absolute conscious- ness, and this direct perception (pratyaksa) unifies the knowable into a single, undivided whole. This is the central concept behind a doctrine originally expounded by Narasimha called 'the non-dualism of direct perception' (pratyaksadvaita). This states that consciousness is essentially perceptive and that its perception of all things operates throughout the universe. Insofar

as phenomena are clearly evident (sphufa) to us, everything is directly perceived by absolute consciousness, with which our individual consciousness is identical. This direct perception unfolds everywhere; the one true reality, it is alone and without companion or rival (nihsapatna). Even though it remains one, it can, by its very nature, perceive distinctions (bheda) between one entity and another, without this engendering any division within it. We distinguish between two entities in empirical terms on the basis of their mutual exclusion (anyonyabhava). The relative distinction {bheda) between them is essentially the perceived difference between their respective characteristics. Despite this difference they are united within the purview of a single cognition insofar as they are equally both manifest appearances. This cognition is the undivided essence (rasa) or 'own nature' (svabhava) of both. Encompassed by the 'fire of consciousness', there is no essential difference between them. Just as when an emerald and ruby reflect each other's light, the ruby is reddish-green and the emerald greenish-red, similarly everything is connected with everything else as part of the single variegated (vicitra) cognition of absolute consciousness. Mahesvarananda writes: The Supreme Lord's unique state of emotivity (asadharanabhava) is the outpouring of pure Being (mahasatta). It is manifest as the brilliance (sphuratta) of the universe which, if we ponder deeply, [is realized to be] the single flavour (ekarasa) of the essence of Beauty which is the vibration of the bliss of one's own nature. In this way all things are in reality one although divided from the one another sharing as they do the 'single flavour' (ekarasa) of the pure vibration of consciousness. Kashmiri Saiva Realism Kashmiri Saivism as a whole has been variously called a form of 'realistic idealism', 'monistic idealism', ‘idealistic monism' and 'concrete monism'. It is easy to understand why Kashmiri Saivism is The Doctrine of Vibration said to be 'idealistic' and 'monistic', but in what sense is it also 'realistic*? The answer to this question is of no small importance in trying to under- stand the central idea behind its metaphysics and the fundamental importance of the concept of Spanda, in this seemingly impossible marriage between monistic idealism and pluralistic realism. The Kashmiri Saiva approach understands the world to be a symbol of the absolute, that is, as the manner in which it presents itself to us. Again we can

contrast this view with that of the Advaita Vedanta. The Advaita Vedanta understands the world to be an expression of the absolute insofar as it exists by virtue of the absolute's Being. Being is understood to be the real unity which underlies empirically manifest separateness and as such is never empirically manifest. It is only transcendentally actual as 'being-in-itself. The Kashmiri Saiva position represents, in a sense, a reversal of this point of view. The nature of the absolute, and also that of Being, is onceived as an eternal becoming (satatodita), a dynamic flux or Spanda, 'the agency of the act of being'. It is dentified with the concrete actuality of the fact of appearing, not passive unmanifest Being. Appearance (abhasa) alone is real.ppearing (prakasamanatva) is equivalent to the fact of being (astitva). Ksemaraja writes in his commentary on the Stanzas on Vibration: Indeed, all things are manifest because they are nothing but manifestation. The point being that nothing is manifest apart from manifestation The absolutely unmanifest, from this point of view, can have as little existence as the space in a lattice window of a sky-palace. Nay, even less, because even that space can appear as an imagined mage manifest within consciousness. Everything is real according to the manner in which it appears. Even an illusion is in this sense real, insofar as it appears and is known in the manner in which it appears. The empirical and the real are identical categories of thought. As Abhinava says: Thus this is the supreme doctrine (upanisatf), namely that, when- ever and in whatever form [an entity] appears, that then is its particular nature. Perhaps at this stage a brief comparison with Heidegger's ideas might prove to be enlightening and not altogether out of place. According to Heidegger's phenomenology of Being, reality is intelligible in a two-fold Integral Monism of Kashmiri Saivism manner as 'phenomenon' and 'logos'. Heidegger defines what he means by 'phenomenon' as: "that-which-shows-itself. The manifest . . . phenomena are then the collection of that which lies open in broad daylight or can be brought to the light of day — what the Greeks at times implicitly identified as 'ta onta' (the things-which-are)".

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb) heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna**

**Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version) **Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb) strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel

period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar (=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI: 10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills (SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear

differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black**

**holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-

relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb)

the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this

version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstate thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11 August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of

entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.** Characterizing the behaviour of strongly coupled

quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at λ=λc, the correlation length diverges as ξ∼|λ−λc|−ν, while (e) the gap between the ground state and first excited state vanishes as (e) ΔE∼ξ−z∼|λ−λc|zν. Exactly at the critical point λ=λc, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as ω∼|k−k0|z. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-**

**equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004 STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element.

**Model has been applied for the following results:**

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| **Searches for Sterile Neutrinos with the IceCube Detector M. G. Aartsen et al. (IceCube Collaboration) Phys.**<br><br>**Rev. Lett 117, 071801 – Published 8 August 2016 Physics logo See Viewpoint: Hunting the Sterile Neutrino** |
| **(1) The IceCube neutrino telescope at the South Pole has measured the atmospheric muon neutrino spectrum as (=) a function of zenith angle and (e&eb) energy in the approximate 320 GeV to 20 TeV range, to search for (e) the oscillation signatures of (e) light sterile neutrinos.** |
| **(2) No evidence for anomalous νμ or ν¯μ disappearance is observed in (eb) either of two independently developed analyses, each using one year of (e) atmospheric neutrino data.** |
| **(3) New exclusion limits are placed on (e&eb) the parameter space of (e) the 3+1 model, in which muon antineutrinos experience (e) a strong Mikheyev-Smirnov-Wolfenstein-resonant oscillation.** |
| **(4) The exclusion limits extend to (e) sin22024≤0.02 at Δm2∼0.3 eV2 at the 90% confidence level.** |
| **(5) The allowed region from (e) global analysis of appearance experiments, including (e) LSND and MiniBooNE, is excluded at approximately the 99% confidence level for (e) the global best-fit value of \|Ue4\|2. Received 5 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.117.071801 © 2016 American Physical Society** |
| **Synopsis: Searching for Majorana Neutrinos August 16, 2016The KamLAND-Zen collaboration has run the most sensitive search to date for a radioactive decay that could reveal whether neutrinos are Majorana fermions. Synopsis figureKamLAND-Zen Collaboration** |
| **(6) Are neutrinos and (=) antineutrinos the same particles, as hypothesized by (e) physicist Ettore Majorana?** |
| **(7) If true, this "Majorana" nature of neutrinos could explain (eb) why neutrinos have (e) mass and why the Universe has (e) more matter than antimatter.** |
| **(8) Using an underground neutrino detector in Japan, the KamLAND-Zen collaboration has carried out the most sensitive search to date for a process called neutrinoless double-beta decay, which could confirm (eb) Majorana's hypothesis.** |
| **(9) The results set stringent constraints on the possible properties of (e) Majorana neutrinos.** |
| **(10) Double-beta decay is (=) a nuclear process in which two neutrons turn into (e&eb) two protons (or vice versa), emitting (eb) two antineutrinos.** |
| **(11) If neutrinos are Majorana fermions, (eb) the two antineutrinos could annihilate each other, and some double-beta decays would emit (eb) no antineutrinos.** |
| **(12) However, neutrinoless decay is believed to be (=) extremely rare. Searches of the process using (e) large amounts of the 35 naturally occurring (eb) isotopes that can undergo (e&eb) double-beta decay (such as xenon-136) have so far come up empty.** |

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| **(13) The KamLAND-Zen collaboration boosted their experimental sensitivity to (e) neutrinoless double-beta decays through (e&eb) a combination of factors.** |
| **(14) The researchers deployed clean detectors with (e&eb) very low background noise.** |
| **(15) And they used an unprecedented amount of xenon-136, which they purified to (e) remove any radioactive contaminants that could produce (eb) unwanted signals in the detectors.** |
| **(16) Combining data from several years, they improved (eb) the limit on the probability of neutrinoless double-beta decay by (e) sixfold compared to previous searches.** |
| **(17) This probability corresponds to (e&eb) a decay lifetime whose staggering value exceeds (e) 1026 years.** |
| **(18) This measurement allowed (eb) the researchers to set (eb) the most stringent upper limit on (e&eb) the possible mass of Majorana neutrinos (the particle must be lighter than 61–165 meV). This research is published in Physical Review Letters. –Matteo Rini Matteo Rini is the Deputy Editor of Physics.** |
| **How to Reach a Thousand-Second in-Plane Polarization Lifetime with 0.97−GeV/c Deuterons in a Storage Ring G. Guidoboni et al. (JEDI Collaboration) Phys. Rev. Lett. 117, 054801 – Published 28 July 2016 Physics logo See Synopsis: Deuterons Spin together for a Thousand Seconds** |
| **(19) Authors observe a deuteron beam polarization lifetime near 1000 s in (eb) the horizontal plane of (e) a magnetic storage ring (COSY).** |
| **(20) This long spin coherence time is maintained through (e&eb) a combination of beam bunching, electron cooling, (e&eb) sextupole field corrections, and (e&eb) the suppression of collective effects through (e&eb) beam current limits.** |
| **(21) This record lifetime is required for (e) a storage ring search for (e) an intrinsic electric dipole moment on (e&eb) the deuteron at a statistical sensitivity level approaching 10−29 e cm.** |
| **Synopsis: Deuterons Spin together for a Thousand Seconds July 28, 2016 Researchers set a new record for the in-plane spin-alignment lifetime of deuterons circulating in a magnetic storage ring. Synopsis figure E. Stephenson/Indiana University Center for Spacetime Symmetries Sime sentences and figures are deleted due to spatial restrictions.** |
| **(22) The electric dipole moment (EDM) of (e) a fundamental particle or a simple nucleus is believed to be (=) extremely small.** |
| **(23) One might therefore think that it isn't worth measuring this quantity. But it is. A nonzero EDM would signal (eb) a broken symmetry of (e) nature known as CP violation, which is required to (e) explain why the Universe has (e) more matter than antimatter.** |
| **(24) Edward Stephenson from Indiana University, Bloomington, and colleagues has now made a significant stride towards determining (eb) the EDM of the deuteron using (e) a deuteron beam circulating in (e&eb) a magnetic storage ring.** |

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| **(25) They have been able to preserve the alignment of the deuterons' spins for (e) a chunk of time that may be long enough to determine (eb) the deuteron's EDM.** |
| **(26) In experiments of this kind, the EDM can be determined by (e) measuring the rotation of (e) the initially aligned spins of (e) the beam particles as they travel around (e&eb) the ring.** |
| **(27) To obtain a measurement sufficiently accurate to explore (e&eb) CP violation predicted by (e) theories beyond the standard model, the spin alignment has to be (=) preserved for about 1000 s.** |
| **(28) But changes in particle speed—associated with (e&eb) position or direction variations among (e&eb) the particles in the beam—lead to (eb) a rapid loss of the alignment in the plane of the ring.** |
| **(29) Stephenson and co-workers have managed to retain (eb)the alignment of deuterons for the required 1000 s, which is 3 orders of magnitude longer (eb) than the previous record (obtained for an electron beam).** |
| **(30) Their success resulted from (e) capturing the beam into a tight bunch to reduce (e) position and direction spread and from (e) shaping the ring's magnetic field to(e) ensure all of the deuteron orbits around (e&eb) the ring have the same length, despite (e) direction variations. This research is published in Physical Review Letters. –Ana Lopes Ana Lopes is a Senior Editor of Physics.** |
| **Cotunneling Drag Effect in Coulomb-Coupled Quantum Dots A. J. Keller, J. S. Lim, David Sánchez, Rosa López, S. Amasha, J. A. Katine, Hadas Shtrikman, and D. Goldhaber-Gordon Phys. Rev. Lett. 117, 066602 – Published 4 August 2016 Physics logo See Synopsis: Coulomb Drag in a Double Dot** |
| **(31) In Coulomb drag, a current flowing in one conductor can induce (eb) a voltage across an adjacent conductor via (e&eb) the Coulomb interaction.** |
| **(32) The mechanisms yielding drag effects are not always understood, even though drag effects are (=) sufficiently general to be (=) seen in many low-dimensional systems.** |
| **(33) In this Letter, authors observe Coulomb drag in (eb) a Coulomb-coupled double quantum dot and, through (e&eb) both experimental and theoretical arguments, identify cotunneling as (=) essential to obtaining a correct qualitative understanding of the drag behavior.** |
| **Synopsis: Coulomb Drag in a Double Dot August 4, 2016 Electric current passing through a quantum dot can generate current in a nearby dot through a coordinated tunneling mechanism. Synopsis figureA. J. Keller et al., Phys. Rev. Lett. (2016)** |
| **(34) Electrons moving through (e&eb) a conductor can induce (eb) the movement of electrons of a nearby, but isolated, conductor.** |
| **(35) This so-called Coulomb drag is well understood in 2D layered materials, but the theory for nanowires and quantum dots does not (e) fully explain experiments.** |
| **(36) A new study of Coulomb drag in two nearby quantum dots shows (eb) that simultaneous tunneling of electrons—called cotunneling—can help explain the low-temperature behavior of (e) the system.** |

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| **(37) Electric current flows through (e&eb) a quantum dot by the tunneling of (e) electrons.** |
| **(38) When two dots are placed next to each other, the active, or lead, dot can drag current through (e&eb) the passive dot by (e) a process called sequential tunneling.** |
| **(39) In this case, an electron first tunnels onto (e&eb) the active dot, increasing (eb) the probability of tunneling onto (e&eb) the passive dot.** |
| **(40) Thus the passive dot follows (e&eb) the active dot's lead, but always one step behind.** |
| **(41) Sequential tunneling explains (eb) Coulomb drag observations at (eb) high temperature, but no studies have investigated cold temperatures where (e) quantum coherence effects may play a role.** |
| **(42) Andrew Keller from the California Institute of Technology, Pasadena, and his colleagues explored Coulomb drag at (eb) sub-Kelvin temperatures in a pair of quantum dots coupled by (e) capacitance.** |
| **(43) The team applied a current-driving voltage across (e&eb) one dot and detected a drag current in (eb) the other dot, even in regimes where (e) the driving voltage was very small.** |
| **(44) This conflicts with (e&eb) the sequential tunneling model, which predicts (eb) a threshold driving voltage below which (e) the drag current should be zero.** |
| **(45) The team showed that the inclusion of cotunneling into (eb) the drag theory explains the observations.** |
| **(46) Cotunneling could be important to efforts aimed at harvesting a unidirectional current from (e) a fluctuating voltage supply. This research is published in Physical Review Letters. –Michael Schirber Michael Schirber is a Corresponding Editor for Physics based in Lyon, France.** |
| **Maybe it wasn't the Higgs particle after all November 7, 2014 Maybe it wasn't the Higgs particle after all Associate Professor Mads Toudal Frandsen, University of Southern Denmark Credit: University of Southern Denmark** |
| **(47) Last year CERN announced the finding of a new elementary particle, the Higgs particle. But maybe it wasn't the Higgs particle, maybe it just looks like it. And maybe it is not alone. Many calculations indicate that the particle discovered last year in the CERN particle accelerator was (=) indeed the famous Higgs particle.** |
| **(48) Physicists agree that the CERN experiments did find a new particle that had never been seen before, but according to (e&eb) an international research team, there is (=) no conclusive evidence that the particle was indeed the Higgs particle.The research team has scrutinized the existing scientific data from CERN about the newfound particle and published their analysis in the journal Physical Review D. A member of this team is Mads Toudal Frandsen, associate professor at the Center for Cosmology and Particle Physics Phenomenology, Department of Physics, Chemistry and Pharmacy at the University of Southern Denmark. "The CERN data is generally taken as (=) evidence that the particle is the Higgs particle.** |

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| **(49) It is true that the Higgs particle can explain (eb) the data but there can be other explanations, we would also get this data from other particles", Mads Toudal Frandsen explains.** |
| **(50) The researchers' analysis does not (e) debunk the possibility that CERN has discovered the Higgs particle. That is still possible - but it is equally possible that it is a different kind of particle. "The current data is not precise enough to determine exactly what the particle is. It could be a number of other known particles", says Mads Toudal Frandsen. But if it wasn't the Higgs particle, which was found in CERN's particle accelerator, then what was it? "We believe that it may be a so-called techni-higgs particle. This particle is in some ways similar to the Higgs particle - hence half of the name", says Mads Toudal Frandsen.** |
| **(51) Although the techni-higgs particle and Higgs particle can easily be confused in experiments, they are two very different particles belonging to (e) two very different theories of how the universe was created.** |
| **(52) The Higgs particle is the missing piece in (eb) the theory called the Standard Model.** |
| **(53) This theory describes (eb) three of the four forces of nature. But it does not (e) explain what dark matter is - the substance that makes up (eb) most of the universe. A techni-higgs particle, if it exists, is (=) a completely different thing:** |
| **(54) "A techni-higgs particle is not (e) an elementary particle.** |
| **(55) Instead, it consists of so-called techni-quarks, which we believe are elementary. Techni-quarks may bind together in various ways to form (eb) for instance techni-higgs particles, while other combinations may form (eb) dark matter. Authors therefore expect to find several different particles at the LHC, all built by (e) techni-quarks", says Mads Toudal Frandsen.** |
| **(56) If techni-quarks exist, (eb) there must be a force to bind them together so that (e) they can form particles.** |
| **(57) None of the four known forces of nature (gravity, the electromagnetic force, the weak nuclear force and the strong nuclear force) are (=) any good at binding techni-quarks together.** |
| **(58) There must therefore be (=) a yet undiscovered force of nature. This force is called the the Technicolor force.** |
| **(59) What was found last year in CERN's accelerator could thus be (=) either the Higgs particle of the Standard Model or a light techni-higgs particle, composed of (e) two techni-quarks.** |
| **(60) Mads Toudal Frandsen believes that more data from CERN will probably be able to determine (eb) if it was a Higgs or a techni-higgs particle. If CERN gets an even more powerful accelerator, it will in principle be able to observe techni-quarks directly. More information: Technicolor Higgs boson in the light of LHC data. Phys. Rev. D 90, 035012th Alexander Belyaev, Matthew S. Brown, Roshan Foadi, and Mads T. Frandsen. journals.aps.org/prd/abstract/10.1103/PhysRevD.90.035012. On Arxiv: arxiv.org/abs/1309.2097 Journal reference: Physical Review D search and more info website: Provided by: University of Southern Denmark search and more info website** |

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| **The inflationary brane-antibrane universeClifford P. Burgess1,2, Mahbub Majumdar3, Detlef Nolte1, Fernando Quevedo3, Govindan Rajesh1 and Ren-Jie Zhang1 Published 13 August 2001 • Journal of High Energy Physics, Volume 2001, JHEP07(2001)** |
| **(61) Authors show how the motion through (e&eb) the extra dimensions of a gas of branes and antibranes can, under certain circumstances, produce (eb) an era of inflation as seen by (e) observers trapped on (e&eb) a 3-brane, with (e&eb) the inflaton being the inter-brane separation.** |
| **(62) Although most of discussion refers to (e) arbitrary ps, when we need to be specific they assume that they are (=) D-branes of type-II or type-I string theory.** |
| **(63) For realistic brane couplings, such as those arising in (eb) string theory, the inter-brane potentials are (=) too steep to inflate (eb) the universe for acceptably long times.** |
| **(64) However, for special regions of the parameter space of (e) brane-antibrane positions the brane motion is (=) slow enough for there to be (=) sufficient inflation.** |
| **(65) Inflation would be more generic in models where (e) the inter-brane interactions are much weaker.** |
| **(66) The spectrum of primordial density fluctuations predicted has (e) index n slightly less than 1, and acceptable amplitude, provided (e) that the extra dimensions have (e) linear size 1/r ~ 1012 GeV.** |
| **(67) Reheating occurs as in hybrid inflation, with (e&eb) the tachyonic instability of (e) the brane-antibrane system taking over for (e) small separations.** |
| **(68) The tachyon field can induce (eb) a cascade mechanism within which higher-dimension branes annihilate (e) into lower-dimension ones.** |
| **(69) They argue that such a cascade naturally stops with (e&eb) the production of 3-branes in (eb) 10-dimensional string theory.** |
| **Effect of the neutron lifetime on processes in the early universe A. V. Chechkin, A. V. Ivanchik, A. P. Serebrov, S. V. Bobashev First Online: 27 July 2016 DOI: 10.1134/S1063784216070069 Chechkin, A.V., Ivanchik, A.V., Serebrov, A.P. et al. Tech. Phys. (2016) 61: 1101. Doi: 10.1134/S1063784216070069** |
| **(70) The influence of the neutron lifetime on (e&eb) the abundance of light elements produced during (e&eb) the primordial nucleosynthesis early in (eb) the birth of the Universe is considered.** |
| **(71) Among light elements, namely, D, 3He, 4He, and 7Li, 4He proves to be (=) most sensitive to neutron lifetime τn.** |
| **(72) Astronomic data on the light element abundance also provide (eb) the best accuracy for 4He.** |
| **(73) The solution of a number of problems discussed in this paper requires (e) improving the accuracy of observations for the 4He abundance and refining the value of τn.** |

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| **Application of a Theory and Simulation based Convective Boundary Mixing model for AGB Star Evolution and Nucleosynthesis U. Battino, M. Pignatari, C. Ritter, F. Herwig, P. Denisenkov, J.W. Den Hartogh, R. Trappitsch, R. Hirschi, B. Freytag, F. Thielemann, B. Paxton (Submitted on 19 May 2016)** |
| **(74) The s-process nucleosynthesis in Asymptotic Giant Branch (AGB) stars depends on (e&eb) the modeling of (e) convective boundaries.** |
| **(75) Authors' present models and s-process simulations that adopt (e&eb) a treatment of convective boundaries based on (e) the results of hydrodynamic simulations and on (e&eb) the theory of mixing due to (e) gravity waves in (eb) the vicinity of convective boundaries.** |
| **(76) Hydrodynamics simulations suggest (eb) the presence of convective boundary mixing (CBM) at (eb) the bottom of the thermal pulse-driven convective zone.** |
| **(77) Similarly, convection-induced mixing processes are proposed for (e) the mixing below the convective envelope during (e&eb) third dredge-up where (e) the 13C pocket for (e) the s process in AGB stars forms.** |
| **(78) In this work apply apply a CBM model motivated by (e) simulations and theory to (e&eb) models with initial mass M = 2 and M = 3M?, and with (e&eb) initial metal content Z = 0:01 and Z = 0:02.** |
| **(79) As reported previously, the He-intershell abundance of 12C and 16O are increased by (e) CBM at the bottom of (e) pulse-driven convection zone.** |
| **(80) This mixing is affecting (e&eb) the 22Ne (alpha, n) 25Mg activation and (e&eb) the s-process effciency in the 13C-pocket.** |
| **(81) In our model CBM at the bottom of the convective envelope during (e&eb) the third dredgeup represents (eb) gravity wave mixing.** |
| **(82) Authors take further into account that hydrodynamic simulations indicate (eb) a declining mixing efficiency already about (e&eb) a pressure scale height from (e) the convective boundaries, compared to (e&eb) mixing-length theory.** |
| **(83) The final s-process abundances are characterized by (e) 0:36 < [s=Fe] < 0:78 and the heavy-to-light s-process ratio is (=) 0:23 < [hs=ls] < 0:45.** |
| **(84) Finally, authors compare results with (e&eb) stellar observations, pre-solar grain measurements and previous work. Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI:10.3847/0004-637X/827/1/30 Cite as: arXiv:1605.06159 [astro-ph.SR] (or arXiv:1605.06159v1 [astro-ph.SR] for this version)** |
| **MESA and NuGrid simulations of classical novae: CO and ONe nova nucleosynthesis Pavel A. Denissenkov (UVic), James W. Truran (Chicago), Marco Pignatari (Basel), Reto Trappitsch (Chicago), Christian Ritter (Frankfurt), Falk Herwig (UVic), Umberto Battino (Basel), Kiana Setoodehnia (McMaster), B. Paxton (UC Santa Barbara)** |

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| **(85) Classical novae are the result of (e) thermonuclear flashes of hydrogen accreted by (e) CO or ONe white dwarfs, leading (eb) eventually to the dynamic ejection of (e) the surface layers.** |
| **(86) These are observationally known to be (=) enriched in heavy elements, such as (=) C, O and Ne that must originate in (eb) layers below the H-flash convection zone.** |
| **(87) Building on their previous work, they now present stellar evolution simulations of (e) ONe novae and provide (eb) a comprehensive comparison of models with (e&eb) published ones.** |
| **(88) Some of authors models include exponential convective boundary mixing to account for (e) the observed enrichment of (e) the nova ejecta even when accreted material has (e) a solar abundance distribution.** |
| **(89) Models produce (eb) maximum temperature evolution profiles and nucleosynthesis yields (eb) in good agreement with (e&eb) models that generate (eb) enriched ejecta by assuming (eb) that the accreted material was pre-mixed.** |
| **(90) Authors confirm for (e) ONe novae the result we reported previously, i.e.\ we found that 3He could be produced {\it in situ} in (eb) solar-composition envelopes accreted with (e&eb) slow rates ($M˙<10-10M⊙$/yr) by (e) cold (TWD<107 K) CO WDs, and that convection was triggered by (eb) 3He burning before (e) the nova outburst in that case.** |
| **(91) In addition, authors now find that the interplay between the 3He production and (e&eb) destruction in the solar-composition envelope accreted with (e&eb) an intermediate rate, e.g.\ $M˙=10-10M⊙$/yr, by the 1.15$M⊙$ ONe WD with (e&eb) a relatively high initial central temperature, e.g.\ TWD=15×106 K, leads to (eb) the formation of a thick radiative buffer zone that separates (e&eb) the bottom of the convective envelope from (e) the WD surface. (Abridged) Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI: 10.1093/mnras/stu1000 Cite as: arXiv:1303.6265 [astro-ph.SR] (or arXiv:1303.6265v2 [astro-ph.SR] for this version)** |
| **Hybrid C-O-Ne White Dwarfs as Progenitors of Diverse SNe Ia Pavel Denissenkov (UVic), James Truran (U. Chicago), Falk Herwig (UVic), Sam Jones (UVic), Bill Paxton (UC Santa Barbara), Ken'ichi Nomoto (U. Tokyo), Toshio Suzuki (Nihon U.), Hiroshi Toki (Osaka U.)(Submitted on 6 Nov 2014)** |
| **(92) When carbon is ignited off-center in a CO core of a super-AGB star (eb) its burning in a convective shell tends to propagate to (e&eb) the center.** |
| **(93) Whether the C flame will actually be able to reach (eb) the center depends on (e&eb) the efficiency of (e) extra mixing beneath (e&eb) the C convective shell.** |
| **(94) Whereas thermohaline mixing is too inefficient to interfere with the C-flame propagation, (eb) convective boundary mixing can prevent (e) the C burning from reaching (eb) the center.** |
| **(95) As a result, a C-O-Ne white dwarf (WD) is (=) formed, after the star has lost (e) its envelope.** |
| **(96) Such a "hybrid" WD has (e) a small CO core surrounded by (e) a thick ONe zone.** |

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| **(97) In the proposed 1D stellar evolution computations the hybrid WD is allowed to (eb) accrete C-rich material, as if it were in a close binary system and (e&eb) accreted H-rich material from (e) its companion with (e&eb) a sufficiently high rate at which the accreted H would be processed into (e&eb) He under (e&eb) stationary conditions, assuming (eb) that He could then be transformed into (e&eb) C.** |
| **(98) When the mass of the accreting WD approaches (eb) the Chandrasekhar limit, they find a series of convective Urca shell flashes associated with (e&eb) high abundances of 23Na and 25Mg.** |
| **(99) They are followed by off-center C ignition leading to (eb) convection that occupies (e) almost the entire star.** |
| **(100) To model the Urca processes, authors use the most recent well-resolved data for (e) their reaction and (e&eb) neutrino-energy loss rates.** |
| **(101) Because of the emphasized uncertainty of the convective Urca process in hybrid WD models of (e) SN Ia progenitors, authors consider a number of their potentially possible alternative instances for (e) different mixing assumptions, all of which reach (eb) a phase of explosive C ignition, either off or in (eb) the center.** |
| **(102) Our hybrid SN Ia progenitor models have (e) much lower C to O abundance ratios at (eb) the moment of the explosive C ignition than (e) their pure CO counterparts, which may explain (eb) the observed diversity of the SNe Ia. subjects: Solar and Stellar Astrophysics (astro-ph.SR)cite as: arXiv:1411.1471 [astro-ph.SR](or arXiv:1411.1471v1 [astro-ph.SR] for this version)** |
| **Rejecting proposed dense-matter equations of state with quiescent low-mass X-ray binaries Sebastien Guillot, Robert E. Rutledge (Submitted on 15 Sep 2014 (v1), last revised 20 Oct 2014 (this version, v2))** |
| **(103) Neutrons stars are unique laboratories to (e) discriminate between the various proposed equations of (e) state of matter at and above nuclear density.** |
| **(104) One sub-class of neutron stars - those inside (eb) quiescent low-mass X-ray binaries (qLMXBs) - produce (eb) a thermal surface emission from (e) which the neutron star radius (R_NS) can be measured, using (e) the widely accepted observational scenario for qLMXBs, assuming (eb) unmagnetized H atmospheres.** |
| **(105) In a combined spectral analysis, this work first reproduces (eb) a previously published measurement of the \rns, assumed to be (=) the same for all neutron stars, using (e) a slightly expanded data set.** |
| **(106) The radius measured is R_NS = 9.4 +/-1.2 km. On the basis of spectral analysis alone, this measured value is not affected by imposing (e&eb) an assumption of causality in the core.** |
| **(107) However, the assumptions underlying this R_NS measurement would be falsified by (e) the observation of any neutron star with (e&eb) a mass >2.6 Msun, since radii <11 km would be rejected if (e) causality is assumed, which would exclude (eb) most of the R_NS parameter space obtained in this analysis.** |

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| **(108) Finally, this work directly tests (e&eb) a selection of dense matter equations of states: WFF1, AP4, MPA1, PAL1, MS0, and three versions of equations of state produced through (e&eb) chiral effective theory.** |
| **(109) Two of those, MSO and PAL1, are rejected at (eb) the 99% confidence level, accounting for (e) all quantifiable uncertainties, while (e) the other cannot be excluded at >99% certainty. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE) DOI: 10.1088/2041-8205/796/1/L3 Cite as: arXiv:1409.4306 [astro-ph.HE] (or arXiv:1409.4306v2 [astro-ph.HE] for this version** |

## SECTION EIGHT

## Stochastic Gravitational-Wave Background From Binary Black Holes

### INTRODUCTION—VARIABLES USED

**GW150914: Implications for the Stochastic Gravitational-Wave Background from Binary Black Holes B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 131102 – Published 31 March 2016**

(1) The LIGO detection of the gravitational wave transient GW150914, from (e) the inspiral and merger of (e) two black holes with masses ≳30M⊙, suggests (eb) a population of binary black holes with (e&eb) relatively high mass.
(2) This observation implies (eb) that the stochastic gravitational-wave background from binary black holes, created from (e) the incoherent superposition of (e) all the merging binaries in the Universe, could be (=) higher than previously expected.
(3) Using the properties of GW150914, authors estimate the energy density of (e) such a background from (e) binary black holes.
(4) In the most sensitive part of the Advanced LIGO and Advanced Virgo band for (e) stochastic backgrounds (near 25 Hz), authors predict $\Omega_{GW}(f=25\ Hz)=1.1^{+2.7}_{-0.9}\times10^{-9}$ with 90% confidence.
(5) This prediction is robustly demonstrated for (e) a variety of formation scenarios with (e&eb) different parameters.
(6) The differences between models are (=) small compared to the statistical uncertainty arising from (e) the currently poorly constrained (e) local coalescence rate.
(7) Authors conclude that this background is potentially measurable by (e) the Advanced LIGO and Advanced Virgo detectors operating at their projected final sensitivity. DOI:http://dx.doi.org/10.1103/PhysRevLett.116.131102 © 2016 American Physical Societ

### NOTATION

**Module One**

LIGO detection of the gravitational wave transient GW150914, from the inspiral and merger of two black holes with masses ≳30M⊙, implies that the stochastic gravitational-wave background from binary black holes, created from (e) the incoherent superposition of (e) all the merging binaries in the Universe, could be

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| (=) higher than previously expected | |
| $G_{13}$ : Category one of incoherent superposition of (e) all the merging binaries in the Universe, could be (=) higher than previously expected<br><br>$G_{14}$ : Category two of SAS<br><br>$G_{15}$ : Category three of SAS | |
| $T_{13}$ : Category one ofLIGO detection of the gravitational wave transient GW150914, from the inspiral and merger of two black holes with masses $\gtrsim 30M\odot$, implies that the stochastic gravitational-wave background from binary black holes, created<br><br>$T_{14}$ : Category two of SAS<br><br>$T_{15}$ : Category three of SAS | |
| **Module Two**<br><br>LIGO detection of the gravitational wave transient GW150914, from the inspiral and merger of two black holes with masses $\gtrsim 30M\odot$, implies that the stochastic gravitational-wave background from binary black holes, created from the incoherent superposition of (e) all the merging binaries in the Universe, could be (=) higher than previously expected | |
| $G_{16}$ : Category one of all the merging binaries in the Universe, could be (=) higher than previously expected<br><br>$G_{17}$ : Category two of SAS<br><br>$G_{18}$ : Category three of SAS | |
| $T_{16}$ : Category one of LIGO detection of the gravitational wave transient GW150914, from the inspiral and merger of two black holes with masses $\gtrsim 30M\odot$, implies that the stochastic gravitational-wave background from binary black holes, created from the incoherent superposition<br><br>$T_{17}$ : Category two of SAS<br><br>$T_{18}$ : Category three of SAS | |
| **Module three**<br><br>LIGO detection of the gravitational wave transient GW150914, from the inspiral and merger of two black holes with masses $\gtrsim 30M\odot$, implies that the stochastic gravitational-wave background from binary black holes, created from the incoherent superposition of all the merging binaries in the Universe, could be (=) higher than previously expected | |
| $G_{20}$ : Category one ofLIGO detection of the gravitational wave transient GW150914, from the inspiral and merger of two black holes with masses $\gtrsim 30M\odot$, implies that the stochastic gravitational-wave background from binary black holes, created from the incoherent superposition of all the merging binaries in the Universe | |

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| $G_{21}$ : Category two of SAS<br>$G_{22}$ : Category three of SAS | |
| $T_{20}$ : Category one of higher than previously expected<br>$T_{21}$ : Category two of SAS<br>$T_{22}$ : Category three of SAS | |
| **Module four**<br>Using the properties of GW150914, authors estimate the energy density of (e) such a background from (e) binary black holes | |
| $G_{24}$ : Category one of such a background from (e) binary black holes<br>$G_{25}$ : Category two of SAS<br>$G_{26}$ : Category three of SAS | |
| $T_{24}$ : Category one of properties of GW150914, authors estimate the energy density<br>$T_{25}$ : Category two of SAS<br>$T_{26}$ : Category three of SAS | |
| **Module five**<br>Using the properties of GW150914, authors estimate the energy density of such a background from (e) binary black holes | |
| $G_{28}$ : Category one of binary black holes<br>$G_{29}$ : Category two of SAS<br>$G_{30}$ : Category three of SAS | |
| $T_{28}$ : Category one of properties of GW150914, authors estimate the energy density of such a background<br>$T_{29}$ : Category two of SAS<br>$T_{30}$ : Category three of SAS | |
| **Module six**<br>In the most sensitive part of the Advanced LIGO and Advanced Virgo band for (e) stochastic backgrounds (near 25 Hz), authors predict ΩGW(f=25 Hz)=1.1+2.7−0.9×10−9 with 90% confidence | |
| $G_{32}$ : Category one of stochastic backgrounds (near 25 Hz), authors predict ΩGW(f=25 Hz)=1.1+2.7−0.9×10−9 with 90% confidence | |

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| $G_{33}$ : Category two of SAS<br>$G_{34}$ : Category three of SAS | |
| $T_{32}$ : Category one of most sensitive part of the Advanced LIGO and Advanced Virgo band<br>$T_{33}$ : Category two of SAS<br>$T_{34}$ : Category three of SAS | |
| **Module seven**<br>$\Omega GW(f=25\ Hz)=1.1+2.7-0.9\times10-9$ | |
| $G_{36}$ : Category one of $\Omega GW(f=25\ Hz)$<br>$G_{37}$ : Category two of SAS<br>$G_{38}$ : Category three of SAS | |
| $T_{36}$ : Category one of$1.1+2.7-0.9\times10-9$ (take all the three categries as equivalent)<br>$T_{37}$ : Category two of SAS<br>$T_{38}$ : Category three of SAS | |
| **Module eight**<br>This prediction is robustly demonstrated for (e) a variety of formation scenarios with (e&eb) different parameters | |
| $G_{40}$ : Category one of variety of formation scenarios with (e&eb) different parameters<br>$G_{41}$ : Category two of SAS<br>$G_{42}$ : Category three of SAS | |
| $T_{40}$ : Category one of prediction is robustly demonstrated<br>$T_{41}$ : Category two of SAS<br>$T_{42}$ : Category three of SAS | |
| **Module Nine**<br>This prediction is robustly demonstrated for a variety of formation scenarios with (e&eb) different parameters | |
| $G_{44}$ : Category one of **prediction is robustly demonstrated for a variety of formation scenarios;** different parameters | |

| | |
|---|---|
| $G_{45}$ : Category two of SAS<br><br>$G_{46}$ : Category three of SAS | |
| $T_{44}$ : Category one of different parameters **;prediction is robustly demonstrated for a variety of formation scenarios**<br><br>$T_{45}$ : Category two of SAS<br><br>$T_{46}$ : Category three of SAS | |
| **The Coefficients:** | |
| $(a_{13})^{(1)}, (a_{14})^{(1)}, (a_{15})^{(1)}, (b_{13})^{(1)}, (b_{14})^{(1)}, (b_{15})^{(1)}\ (a_{16})^{(2)}, (a_{17})^{(2)}, (a_{18})^{(2)}\ (b_{16})^{(2)}, (b_{17})^{(2)}, (b_{18})^{(2)}$:<br>$(a_{20})^{(3)}, (a_{21})^{(3)}, (a_{22})^{(3)}$ , $(b_{20})^{(3)}, (b_{21})^{(3)}, (b_{22})^{(3)}$<br>$(a_{24})^{(4)}, (a_{25})^{(4)}, (a_{26})^{(4)}, (b_{24})^{(4)}, (b_{25})^{(4)}, (b_{26})^{(4)}, (b_{28})^{(5)}, (b_{29})^{(5)}, (b_{30})^{(5)}, (a_{28})^{(5)}, (a_{29})^{(5)}, (a_{30})^{(5)}$,<br>$(a_{32})^{(6)}, (a_{33})^{(6)}, (a_{34})^{(6)}, (b_{32})^{(6)}, (b_{33})^{(6)}, (b_{34})^{(6)}$<br>$(a_{36})^{(7)}, (a_{37})^{(7)}, (a_{38})^{(7)}, (b_{36})^{(7)}, (b_{37})^{(7)}, (b_{38})^{(7)}$<br>$(a_{40})^{(8)}, (a_{41})^{(8)}, (a_{42})^{(8)}, (b_{40})^{(8)}, (b_{41})^{(8)}, (b_{42})^{(8)}$<br>$(a_{44})^{(9)}, (a_{45})^{(9)}, (a_{46})^{(9)}, (b_{44})^{(9)}, (b_{45})^{(9)}, (b_{46})^{(9)}$<br><br>are Accentuation coefficients<br>$(a_{13}^{'})^{(1)}, (a_{14}^{'})^{(1)}, (a_{15}^{'})^{(1)}, (b_{13}^{'})^{(1)}, (b_{14}^{'})^{(1)}, (b_{15}^{'})^{(1)}, (a_{16}^{'})^{(2)}, (a_{17}^{'})^{(2)}, (a_{18}^{'})^{(2)}$, $(b_{16}^{'})^{(2)}, (b_{17}^{'})^{(2)}, (b_{18}^{'})^{(2)}$<br>$, (a_{20}^{'})^{(3)}, (a_{21}^{'})^{(3)}, (a_{22}^{'})^{(3)}, (b_{20}^{'})^{(3)}, (b_{21}^{'})^{(3)}, (b_{22}^{'})^{(3)}$<br>$(a_{24}^{'})^{(4)}, (a_{25}^{'})^{(4)}, (a_{26}^{'})^{(4)}, (b_{24}^{'})^{(4)}, (b_{25}^{'})^{(4)}, (b_{26}^{'})^{(4)}, (b_{28}^{'})^{(5)}, (b_{29}^{'})^{(5)}, (b_{30}^{'})^{(5)}\ (a_{28}^{'})^{(5)}, (a_{29}^{'})^{(5)}, (a_{30}^{'})^{(5)}$ ,<br>$(a_{32}^{'})^{(6)}, (a_{33}^{'})^{(6)}, (a_{34}^{'})^{(6)}, (b_{32}^{'})^{(6)}, (b_{33}^{'})^{(6)}, (b_{34}^{'})^{(6)}$<br>$(a_{36}^{'})^{(7)}, (a_{37}^{'})^{(7)}, (a_{38}^{'})^{(7)}, (b_{36}^{'})^{(7)}, (b_{37}^{'})^{(7)}, (b_{38}^{'})^{(7)}$,<br>$(a_{40}^{'})^{(8)}, (a_{41}^{'})^{(8)}, (a_{42}^{'})^{(8)}, (b_{40}^{'})^{(8)}, (b_{41}^{'})^{(8)}, (b_{42}^{'})^{(8)}$,<br>$(a_{44}^{'})^{(9)}, (a_{45}^{'})^{(9)}, (a_{46}^{'})^{(9)}, (b_{44}^{'})^{(9)}, (b_{45}^{'})^{(9)}, (b_{46}^{'})^{(9)}$,<br>are Dissipation coefficients | |
| **Module Numbered One**<br><br>**The differential system of this model is now (Module Numbered one)** | |
| $\frac{dG_{13}}{dt} = (a_{13})^{(1)}G_{14} - \left[(a_{13}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14}, t)\right]G_{13}$ | 1 |
| $\frac{dG_{14}}{dt} = (a_{14})^{(1)}G_{13} - \left[(a_{14}^{'})^{(1)} + (a_{14}^{''})^{(1)}(T_{14}, t)\right]G_{14}$ | 2 |
| $\frac{dG_{15}}{dt} = (a_{15})^{(1)}G_{14} - \left[(a_{15}^{'})^{(1)} + (a_{15}^{''})^{(1)}(T_{14}, t)\right]G_{15}$ | 3 |
| $\frac{dT_{13}}{dt} = (b_{13})^{(1)}T_{14} - \left[(b_{13}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G, t)\right]T_{13}$ | 4 |
| $\frac{dT_{14}}{dt} = (b_{14})^{(1)}T_{13} - \left[(b_{14}^{'})^{(1)} - (b_{14}^{''})^{(1)}(G, t)\right]T_{14}$ | 5 |
| $\frac{dT_{15}}{dt} = (b_{15})^{(1)}T_{14} - \left[(b_{15}^{'})^{(1)} - (b_{15}^{''})^{(1)}(G, t)\right]T_{15}$ | 6 |
| $+(a_{13}^{''})^{(1)}(T_{14}, t) =$ First augmentation factor | |
| $-(b_{13}^{''})^{(1)}(G, t) =$ First detritions factor | |
| **Module Numbered Two** | |

| The differential system of this model is now ( Module numbered two) | |
|---|---|
| $\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \left[(a_{16}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17},t)\right]G_{16}$ | 7 |
| $\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \left[(a_{17}^{'})^{(2)} + (a_{17}^{''})^{(2)}(T_{17},t)\right]G_{17}$ | 8 |
| $\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \left[(a_{18}^{'})^{(2)} + (a_{18}^{''})^{(2)}(T_{17},t)\right]G_{18}$ | 9 |
| $\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \left[(b_{16}^{'})^{(2)} - (b_{16}^{''})^{(2)}((G_{19}),t)\right]T_{16}$ | 10 |
| $\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \left[(b_{17}^{'})^{(2)} - (b_{17}^{''})^{(2)}((G_{19}),t)\right]T_{17}$ | 11 |
| $\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \left[(b_{18}^{'})^{(2)} - (b_{18}^{''})^{(2)}((G_{19}),t)\right]T_{18}$ | 12 |
| $+(a_{16}^{''})^{(2)}(T_{17},t) =$ First augmentation factor | |
| $-(b_{16}^{''})^{(2)}((G_{19}),t) =$ First detritions factor | |
| **Module Numbered Three**<br><br>**The differential system of this model is now (Module numbered three)** | |
| $\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \left[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21},t)\right]G_{20}$ | 13 |
| $\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \left[(a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)}(T_{21},t)\right]G_{21}$ | 14 |
| $\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \left[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21},t)\right]G_{22}$ | 15 |
| $\frac{dT_{20}}{dt} = (b_{20})^{(3)}T_{21} - \left[(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23},t)\right]T_{20}$ | 16 |
| $\frac{dT_{21}}{dt} = (b_{21})^{(3)}T_{20} - \left[(b_{21}^{'})^{(3)} - (b_{21}^{''})^{(3)}(G_{23},t)\right]T_{21}$ | 17 |
| $\frac{dT_{22}}{dt} = (b_{22})^{(3)}T_{21} - \left[(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23},t)\right]T_{22}$ | 18 |
| $+(a_{20}^{''})^{(3)}(T_{21},t) =$ First augmentation factor | |
| $-(b_{20}^{''})^{(3)}(G_{23},t) =$ First detritions factor | |
| **Module Numbered Four**<br><br>**The differential system of this model is now (Module numbered Four)** | |
| $\frac{dG_{24}}{dt} = (a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25},t)\right]G_{24}$ | 19 |
| $\frac{dG_{25}}{dt} = (a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25},t)\right]G_{25}$ | 20 |
| $\frac{dG_{26}}{dt} = (a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25},t)\right]G_{26}$ | 21 |
| $\frac{dT_{24}}{dt} = (b_{24})^{(4)}T_{25} - \left[(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}((G_{27}),t)\right]T_{24}$ | 22 |
| $\frac{dT_{25}}{dt} = (b_{25})^{(4)}T_{24} - \left[(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}((G_{27}),t)\right]T_{25}$ | 23 |
| $\frac{dT_{26}}{dt} = (b_{26})^{(4)}T_{25} - \left[(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}((G_{27}),t)\right]T_{26}$ | 24 |
| $+(a_{24}^{''})^{(4)}(T_{25},t) =$ **First augmentation factor** | |
| $-(b_{24}^{''})^{(4)}((G_{27}),t) =$ **First detritions factor** | |
| **Module Numbered Five**:<br><br>**The differential system of this model is now (Module number five)** | |
| $\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29},t)\right]G_{28}$ | 25 |
| $\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29},t)\right]G_{29}$ | 26 |

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| $\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \left[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29}, t)\right]G_{30}$ | 27 |
| $\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \left[(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}((G_{31}), t)\right]T_{28}$ | 28 |
| $\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \left[(b_{29}^{'})^{(5)} - (b_{29}^{''})^{(5)}((G_{31}), t)\right]T_{29}$ | 29 |
| $\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \left[(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}((G_{31}), t)\right]T_{30}$ | 30 |
| $+(a_{28}^{''})^{(5)}(T_{29}, t) =$ **First augmentation factor** | |
| $-(b_{28}^{''})^{(5)}((G_{31}), t) =$ **First detritions factor** | |
| **Module Numbered Six**<br><br>**The differential system of this model is now (Module numbered Six)** | |
| $\frac{dG_{32}}{dt} = (a_{32})^{(6)}G_{33} - \left[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33}, t)\right]G_{32}$ | 31 |
| $\frac{dG_{33}}{dt} = (a_{33})^{(6)}G_{32} - \left[(a_{33}^{'})^{(6)} + (a_{33}^{''})^{(6)}(T_{33}, t)\right]G_{33}$ | 32 |
| $\frac{dG_{34}}{dt} = (a_{34})^{(6)}G_{33} - \left[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33}, t)\right]G_{34}$ | 33 |
| $\frac{dT_{32}}{dt} = (b_{32})^{(6)}T_{33} - \left[(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}((G_{35}), t)\right]T_{32}$ | 34 |
| $\frac{dT_{33}}{dt} = (b_{33})^{(6)}T_{32} - \left[(b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}((G_{35}), t)\right]T_{33}$ | 35 |
| $\frac{dT_{34}}{dt} = (b_{34})^{(6)}T_{33} - \left[(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}((G_{35}), t)\right]T_{34}$ | 36 |
| $+(a_{32}^{''})^{(6)}(T_{33}, t) =$ **First augmentation factor** | |
| **Module Numbered Seven:**<br><br>**The differential system of this model is now (Seventh Module)** | |
| $\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \left[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37}, t)\right]G_{36}$ | 37 |
| $\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \left[(a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37}, t)\right]G_{37}$ | *38* |
| $\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \left[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37}, t)\right]G_{38}$ | *39* |
| $\frac{dT_{36}}{dt} = (b_{36})^{(7)}T_{37} - \left[(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}((G_{39}), t)\right]T_{36}$ | *40* |
| $\frac{dT_{37}}{dt} = (b_{37})^{(7)}T_{36} - \left[(b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}((G_{39}), t)\right]T_{37}$ | *41* |
| $\frac{dT_{38}}{dt} = (b_{38})^{(7)}T_{37} - \left[(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}((G_{39}), t)\right]T_{38}$ | *42* |
| $+(a_{36}^{''})^{(7)}(T_{37}, t) =$ **First augmentation factor** | |
| **Module Numbered Eight**<br><br>The differential system of this model is now | |
| $\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \left[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41}, t)\right]G_{40}$ | *43* |
| $\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \left[(a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41}, t)\right]G_{41}$ | *44* |
| $\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \left[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41}, t)\right]G_{42}$ | *45* |
| $\frac{dT_{40}}{dt} = (b_{40})^{(8)}T_{41} - \left[(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}((G_{43}), t)\right]T_{40}$ | *46* |
| $\frac{dT_{41}}{dt} = (b_{41})^{(8)}T_{40} - \left[(b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}((G_{43}), t)\right]T_{41}$ | *47* |
| $\frac{dT_{42}}{dt} = (b_{42})^{(8)}T_{41} - \left[(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}((G_{43}), t)\right]T_{42}$ | *48* |
| **Module Numbered Nine** | |

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| The differential system of this model is now | |
| $\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \left[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45},t)\right]G_{44}$ | 49 |
| $\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \left[(a_{45}^{'})^{(9)} + (a_{45}^{''})^{(9)}(T_{45},t)\right]G_{45}$ | 50 |
| $\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \left[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45},t)\right]G_{46}$ | 51 |
| $\frac{dT_{44}}{dt} = (b_{44})^{(9)}T_{45} - \left[(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}((G_{47}),t)\right]T_{44}$ | 52 |
| $\frac{dT_{45}}{dt} = (b_{45})^{(9)}T_{44} - \left[(b_{45}^{'})^{(9)} - (b_{45}^{''})^{(9)}((G_{47}),t)\right]T_{45}$ | 53 |
| $\frac{dT_{46}}{dt} = (b_{46})^{(9)}T_{45} - \left[(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}((G_{47}),t)\right]T_{46}$ | 54 |
| $+(a_{44}^{''})^{(9)}(T_{45},t) =$ **First augmentation factor** | |
| $-(b_{44}^{''})^{(9)}((G_{47}),t) =$ **First detrition factor** | |
| $\frac{dG_{13}}{dt} = (a_{13})^{(1)}G_{14} - \begin{bmatrix} (a_{13}^{'})^{(1)} & +(a_{13}^{''})^{(1)}(T_{14},t) & +(a_{16}^{''})^{(2,2,)}(T_{17},t) & +(a_{20}^{''})^{(3,3,)}(T_{21},t) \\ +(a_{24}^{''})^{(4,4,4,4,)}(T_{25},t) & +(a_{28}^{''})^{(5,5,5,5,)}(T_{29},t) & +(a_{32}^{''})^{(6,6,6,6,)}(T_{33},t) & \\ +(a_{36}^{''})^{(7,7)}(T_{37},t) & +(a_{40}^{''})^{(8,8)}(T_{41},t) & +(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t) & \end{bmatrix} G_{13}$ | *55* |
| $\frac{dG_{14}}{dt} = (a_{14})^{(1)}G_{13} - \begin{bmatrix} (a_{14}^{'})^{(1)} & +(a_{14}^{''})^{(1)}(T_{14},t) & +(a_{17}^{''})^{(2,2,)}(T_{17},t) & +(a_{21}^{''})^{(3,3,)}(T_{21},t) \\ +(a_{25}^{''})^{(4,4,4,4,)}(T_{25},t) & +(a_{29}^{''})^{(5,5,5,5,)}(T_{29},t) & +(a_{33}^{''})^{(6,6,6,6,)}(T_{33},t) & \\ +(a_{37}^{''})^{(7,7)}(T_{37},t) & +(a_{41}^{''})^{(8,8)}(T_{41},t) & +(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t) & \end{bmatrix} G_{14}$ | *56* |
| $\frac{dG_{15}}{dt} = (a_{15})^{(1)}G_{14} - \begin{bmatrix} (a_{15}^{'})^{(1)} & +(a_{15}^{''})^{(1)}(T_{14},t) & +(a_{18}^{''})^{(2,2,)}(T_{17},t) & +(a_{22}^{''})^{(3,3,)}(T_{21},t) \\ +(a_{26}^{''})^{(4,4,4,4,)}(T_{25},t) & +(a_{30}^{''})^{(5,5,5,5,)}(T_{29},t) & +(a_{34}^{''})^{(6,6,6,6,)}(T_{33},t) & \\ +(a_{38}^{''})^{(7,7)}(T_{37},t) & +(a_{42}^{''})^{(8,8)}(T_{41},t) & +(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t) & \end{bmatrix} G_{15}$ | *57* |
| Where $(a_{13}^{''})^{(1)}(T_{14},t)$, $(a_{14}^{''})^{(1)}(T_{14},t)$, $(a_{15}^{''})^{(1)}(T_{14},t)$ are first augmentation coefficients for category 1, 2 and 3<br>$+(a_{16}^{''})^{(2,2,)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,)}(T_{17},t)$ are second augmentation coefficient for category 1, 2 and 3<br>$+(a_{20}^{''})^{(3,3,)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3<br>$+(a_{24}^{''})^{(4,4,4,4,)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3<br>$+(a_{28}^{''})^{(5,5,5,5,)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,)}(T_{29},t)$ are fifth augmentation coefficient for category 1, 2 and 3<br>$+(a_{32}^{''})^{(6,6,6,6,)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,)}(T_{33},t)$ are sixth augmentation coefficient for category 1, 2 and 3<br>$+(a_{38}^{''})^{(7,7)}(T_{37},t)$ $+(a_{37}^{''})^{(7,7)}(T_{37},t)$ $+(a_{36}^{''})^{(7,7)}(T_{37},t)$ are seventh augmentation coefficient for 1,2,3<br>$+(a_{40}^{''})^{(8,8)}(T_{41},t)$ $+(a_{41}^{''})^{(8,8)}(T_{41},t)$ $+(a_{42}^{''})^{(8,8)}(T_{41},t)$ are eight augmentation coefficient for 1,2,3<br>$+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficient for 1,2,3 | |

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| $\frac{dT_{13}}{dt} = (b_{13})^{(1)}T_{14} - \begin{bmatrix} (b_{13}^{'})^{(1)}\boxed{-(b_{13}^{''})^{(1)}(G,t)} & \boxed{-(b_{16}^{''})^{(2,2,)}(G_{19},t)} & \boxed{-(b_{20}^{''})^{(3,3,)}(G_{23},t)} \\ \boxed{-(b_{24}^{''})^{(4,4,4,4,)}(G_{27},t)} & \boxed{-(b_{28}^{''})^{(5,5,5,5,)}(G_{31},t)} & \boxed{-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)} \\ \boxed{-(b_{36}^{''})^{(7,7,)}(G_{39},t)} & \boxed{-(b_{40}^{''})^{(8,8)}(G_{43},t)} & \boxed{-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{13}$ | 58 |
| $\frac{dT_{14}}{dt} = (b_{14})^{(1)}T_{13} - \begin{bmatrix} (b_{14}^{'})^{(1)}\boxed{-(b_{14}^{''})^{(1)}(G,t)} & \boxed{-(b_{17}^{''})^{(2,2,)}(G_{19},t)} & \boxed{-(b_{21}^{''})^{(3,3,)}(G_{23},t)} \\ \boxed{-(b_{25}^{''})^{(4,4,4,4,)}(G_{27},t)} & \boxed{-(b_{29}^{''})^{(5,5,5,5,)}(G_{31},t)} & \boxed{-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)} \\ \boxed{-(b_{37}^{''})^{(7,7,)}(G_{39},t)} & \boxed{-(b_{41}^{''})^{(8,8)}(G_{43},t)} & \boxed{-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{14}$ | 59 |
| $\frac{dT_{15}}{dt} = (b_{15})^{(1)}T_{14} - \begin{bmatrix} (b_{15}^{'})^{(1)}\boxed{-(b_{15}^{''})^{(1)}(G,t)} & \boxed{-(b_{18}^{''})^{(2,2,)}(G_{19},t)} & \boxed{-(b_{22}^{''})^{(3,3,)}(G_{23},t)} \\ \boxed{-(b_{26}^{''})^{(4,4,4,4,)}(G_{27},t)} & \boxed{-(b_{30}^{''})^{(5,5,5,5,)}(G_{31},t)} & \boxed{-(b_{34}^{''})^{(6,6,6,6,)}(G_{35},t)} \\ \boxed{-(b_{38}^{''})^{(7,7,)}(G_{39},t)} & \boxed{-(b_{42}^{''})^{(8,8)}(G_{43},t)} & \boxed{-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{15}$ | 60 |
| Where $\boxed{-(b_{13}^{''})^{(1)}(G,t)}$, $\boxed{-(b_{14}^{''})^{(1)}(G,t)}$, $\boxed{-(b_{15}^{''})^{(1)}(G,t)}$ are first detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{16}^{''})^{(2,2,)}(G_{19},t)}$, $\boxed{-(b_{17}^{''})^{(2,2,)}(G_{19},t)}$, $\boxed{-(b_{18}^{''})^{(2,2,)}(G_{19},t)}$ are second detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{20}^{''})^{(3,3,)}(G_{23},t)}$, $\boxed{-(b_{21}^{''})^{(3,3,)}(G_{23},t)}$, $\boxed{-(b_{22}^{''})^{(3,3,)}(G_{23},t)}$ are third detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{24}^{''})^{(4,4,4,4,)}(G_{27},t)}$, $\boxed{-(b_{25}^{''})^{(4,4,4,4,)}(G_{27},t)}$, $\boxed{-(b_{26}^{''})^{(4,4,4,4,)}(G_{27},t)}$ are fourth detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{28}^{''})^{(5,5,5,5,)}(G_{31},t)}$, $\boxed{-(b_{29}^{''})^{(5,5,5,5,)}(G_{31},t)}$, $\boxed{-(b_{30}^{''})^{(5,5,5,5,)}(G_{31},t)}$ are fifth detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)}$, $\boxed{-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)}$, $\boxed{-(b_{34}^{''})^{(6,6,6,6,)}(G_{35},t)}$ are sixth detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{37}^{''})^{(7,7,)}(G_{39},t)}$, $\boxed{-(b_{36}^{''})^{(7,7,)}(G_{39},t)}$, $\boxed{-(b_{38}^{''})^{(7,7,)}(G_{39},t)}$ are seventh detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{40}^{''})^{(8,8)}(G_{43},t)}$ $\boxed{-(b_{41}^{''})^{(8,8)}(G_{43},t)}$ $\boxed{-(b_{42}^{''})^{(8,8)}(G_{43},t)}$ are eight detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)}$, $\boxed{-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2 and 3 | |
| | |
| $\frac{dG_{16}}{dt} = (a_{16})^{(2)}G_{17} - \begin{bmatrix} (a_{16}^{'})^{(2)}\boxed{+(a_{16}^{''})^{(2)}(T_{17},t)} & \boxed{+(a_{13}^{''})^{(1,1,)}(T_{14},t)} & \boxed{+(a_{20}^{''})^{(3,3,3)}(T_{21},t)} \\ \boxed{+(a_{24}^{''})^{(4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{28}^{''})^{(5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{32}^{''})^{(6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{36}^{''})^{(7,7,7)}(T_{37},t)} & \boxed{+(a_{40}^{''})^{(8,8,8)}(T_{41},t)} & \boxed{+(a_{44}^{''})^{(9,9)}(T_{45},t)} \end{bmatrix} G_{16}$ | 61 |
| $\frac{dG_{17}}{dt} = (a_{17})^{(2)}G_{16} - \begin{bmatrix} (a_{17}^{'})^{(2)}\boxed{+(a_{17}^{''})^{(2)}(T_{17},t)} & \boxed{+(a_{14}^{''})^{(1,1,)}(T_{14},t)} & \boxed{+(a_{21}^{''})^{(3,3,3)}(T_{21},t)} \\ \boxed{+(a_{25}^{''})^{(4,4,4,4,4)}(T_{25},t)} & \boxed{+(a_{29}^{''})^{(5,5,5,5,5)}(T_{29},t)} & \boxed{+(a_{33}^{''})^{(6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a_{37}^{''})^{(7,7,7)}(T_{37},t)} & \boxed{+(a_{41}^{''})^{(8,8,8)}(T_{41},t)} & \boxed{+(a_{45}^{''})^{(9,9)}(T_{45},t)} \end{bmatrix} G_{17}$ | 62 |

$$\frac{dG_{18}}{dt} = (a_{18})^{(2)}G_{17} - \begin{bmatrix} (a'_{18})^{(2)}\boxed{+(a''_{18})^{(2)}(T_{17},t)} & \boxed{+(a''_{15})^{(1,1,)}(T_{14},t)} & \boxed{+(a''_{22})^{(3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{26})^{(4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{30})^{(5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{34})^{(6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{38})^{(7,7,7)}(T_{37},t)} & \boxed{+(a''_{42})^{(8,8,8)}(T_{41},t)} & \boxed{+(a''_{46})^{(9,9)}(T_{45},t)} \end{bmatrix} G_{18} \quad (63)$$

Where $\boxed{+(a''_{16})^{(2)}(T_{17},t)}$, $\boxed{+(a''_{17})^{(2)}(T_{17},t)}$, $\boxed{+(a''_{18})^{(2)}(T_{17},t)}$ are first augmentation coefficients for category 1, 2 and 3

$\boxed{+(a''_{13})^{(1,1,)}(T_{14},t)}$, $\boxed{+(a''_{14})^{(1,1,)}(T_{14},t)}$, $\boxed{+(a''_{15})^{(1,1,)}(T_{14},t)}$ are second augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{20})^{(3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{21})^{(3,3,3)}(T_{21},t)}$, $\boxed{+(a''_{22})^{(3,3,3)}(T_{21},t)}$ are third augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{24})^{(4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{25})^{(4,4,4,4,4)}(T_{25},t)}$, $\boxed{+(a''_{26})^{(4,4,4,4,4)}(T_{25},t)}$ are fourth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{28})^{(5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{29})^{(5,5,5,5,5)}(T_{29},t)}$, $\boxed{+(a''_{30})^{(5,5,5,5,5)}(T_{29},t)}$ are fifth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{32})^{(6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{33})^{(6,6,6,6,6)}(T_{33},t)}$, $\boxed{+(a''_{34})^{(6,6,6,6,6)}(T_{33},t)}$ are sixth augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{36})^{(7,7,7)}(T_{37},t)}$, $\boxed{+(a''_{37})^{(7,7,7)}(T_{37},t)}$, $\boxed{+(a''_{38})^{(7,7,7)}(T_{37},t)}$ are seventh augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{40})^{(8,8,8)}(T_{41},t)}$, $\boxed{+(a''_{41})^{(8,8,8)}(T_{41},t)}$, $\boxed{+(a''_{42})^{(8,8,8)}(T_{41},t)}$ are eight augmentation coefficient for category 1, 2 and 3

$\boxed{+(a''_{44})^{(9,9)}(T_{45},t)}$, $\boxed{+(a''_{45})^{(9,9)}(T_{45},t)}$, $\boxed{+(a''_{46})^{(9,9)}(T_{45},t)}$ are ninth augmentation coefficient for category 1, 2 and 3

$$\frac{dT_{16}}{dt} = (b_{16})^{(2)}T_{17} - \begin{bmatrix} (b'_{16})^{(2)}\boxed{-(b''_{16})^{(2)}(G_{19},t)} & \boxed{-(b''_{13})^{(1,1,)}(G,t)} & \boxed{-(b''_{20})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b''_{24})^{(4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{28})^{(5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{32})^{(6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{36})^{(7,7,7)}(G_{39},t)} & \boxed{-(b''_{40})^{(8,8,8)}(G_{43},t)} & \boxed{-(b''_{44})^{(9,9)}(G_{47},t)} \end{bmatrix} T_{16} \quad (64)$$

$$\frac{dT_{17}}{dt} = (b_{17})^{(2)}T_{16} - \begin{bmatrix} (b'_{17})^{(2)}\boxed{-(b''_{17})^{(2)}(G_{19},t)} & \boxed{-(b''_{14})^{(1,1,)}(G,t)} & \boxed{-(b''_{21})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b''_{25})^{(4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{29})^{(5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{33})^{(6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{37})^{(7,7,7)}(G_{39},t)} & \boxed{-(b''_{41})^{(8,8,8)}(G_{43},t)} & \boxed{-(b''_{45})^{(9,9)}(G_{47},t)} \end{bmatrix} T_{17} \quad (65)$$

$$\frac{dT_{18}}{dt} = (b_{18})^{(2)}T_{17} - \begin{bmatrix} (b'_{18})^{(2)}\boxed{-(b''_{18})^{(2)}(G_{19},t)} & \boxed{-(b''_{15})^{(1,1,)}(G,t)} & \boxed{-(b''_{22})^{(3,3,3,)}(G_{23},t)} \\ \boxed{-(b''_{26})^{(4,4,4,4,4)}(G_{27},t)} & \boxed{-(b''_{30})^{(5,5,5,5,5)}(G_{31},t)} & \boxed{-(b''_{34})^{(6,6,6,6,6)}(G_{35},t)} \\ \boxed{-(b''_{38})^{(7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9)}(G_{47},t)} \end{bmatrix} T_{18} \quad (66)$$

where $\boxed{-(b''_{16})^{(2)}(G_{19},t)}$, $\boxed{-(b''_{17})^{(2)}(G_{19},t)}$, $\boxed{-(b''_{18})^{(2)}(G_{19},t)}$ are first detrition coefficients for category 1, 2 and 3

$\boxed{-(b''_{13})^{(1,1,)}(G,t)}$, $\boxed{-(b''_{14})^{(1,1,)}(G,t)}$, $\boxed{-(b''_{15})^{(1,1,)}(G,t)}$ are second detrition coefficients for category 1,2 and 3

$\boxed{-(b''_{20})^{(3,3,3,)}(G_{23},t)}$, $\boxed{-(b''_{21})^{(3,3,3,)}(G_{23},t)}$, $\boxed{-(b''_{22})^{(3,3,3,)}(G_{23},t)}$ are third detrition coefficients for category 1,2 and 3

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| $-(b_{24}^{''})^{(4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1,2 and 3<br>$-(b_{28}^{''})^{(5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1,2 and 3<br>$-(b_{32}^{''})^{(6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6)}(G_{35},t)$ are sixth detrition coefficients for category 1,2 and 3<br>$-(b_{36}^{''})^{(7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1,2 and 3<br>$-(b_{40}^{''})^{(8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8)}(G_{43},t)$ are eight detrition coefficients for category 1,2 and 3<br>$-(b_{44}^{''})^{(9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1,2 and 3 | |
| $\frac{dG_{20}}{dt} = (a_{20})^{(3)}G_{21} - \left[ \begin{array}{c} (a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21},t)\ +(a_{16}^{''})^{(2,2,2)}(T_{17},t)\ +(a_{13}^{''})^{(1,1,1,)}(T_{14},t) \\ +(a_{24}^{''})^{(4,4,4,4,4,4)}(T_{25},t)\ +(a_{28}^{''})^{(5,5,5,5,5,5)}(T_{29},t)\ +(a_{32}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ +(a_{36}^{''})^{(7,7,7,7)}(T_{37},t)\ +(a_{40}^{''})^{(8,8,8,8)}(T_{41},t)\ +(a_{44}^{''})^{(9,9,9)}(T_{45},t) \end{array} \right] G_{20}$ | 67 |
| $\frac{dG_{21}}{dt} = (a_{21})^{(3)}G_{20} - \left[ \begin{array}{c} (a_{21}^{'})^{(3)}+(a_{21}^{''})^{(3)}(T_{21},t)\ +(a_{17}^{''})^{(2,2,2)}(T_{17},t)\ +(a_{14}^{''})^{(1,1,1,)}(T_{14},t) \\ +(a_{25}^{''})^{(4,4,4,4,4,4)}(T_{25},t)\ +(a_{29}^{''})^{(5,5,5,5,5,5)}(T_{29},t)\ +(a_{33}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ +(a_{37}^{''})^{(7,7,7,7)}(T_{37},t)\ +(a_{41}^{''})^{(8,8,8,8)}(T_{41},t)\ +(a_{45}^{''})^{(9,9,9)}(T_{45},t) \end{array} \right] G_{21}$ | 68 |
| $\frac{dG_{22}}{dt} = (a_{22})^{(3)}G_{21} - \left[ \begin{array}{c} (a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21},t)\ +(a_{18}^{''})^{(2,2,2)}(T_{17},t)\ +(a_{15}^{''})^{(1,1,1,)}(T_{14},t) \\ +(a_{26}^{''})^{(4,4,4,4,4,4)}(T_{25},t)\ +(a_{30}^{''})^{(5,5,5,5,5,5)}(T_{29},t)\ +(a_{34}^{''})^{(6,6,6,6,6,6)}(T_{33},t) \\ +(a_{38}^{''})^{(7,7,7,7)}(T_{37},t)\ +(a_{42}^{''})^{(8,8,8,8)}(T_{41},t)\ +(a_{46}^{''})^{(9,9,9)}(T_{45},t) \end{array} \right] G_{22}$ | 69 |
| $+(a_{20}^{''})^{(3)}(T_{21},t)$, $+(a_{21}^{''})^{(3)}(T_{21},t)$, $+(a_{22}^{''})^{(3)}(T_{21},t)$ are first augmentation coefficients for category 1, 2 and 3<br>$+(a_{16}^{''})^{(2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2)}(T_{17},t)$ are second augmentation coefficients for category 1, 2 and 3<br>$+(a_{13}^{''})^{(1,1,1,)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,)}(T_{14},t)$ are third augmentation coefficients for category 1, 2 and 3<br>$+(a_{24}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficients for category 1, 2 and 3<br>$+(a_{28}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficients for category 1, 2 and 3<br>$+(a_{32}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation coefficients for category 1, 2 and 3<br>$+(a_{36}^{''})^{(7,7,7,7)}(T_{37},t)$, $+(a_{37}^{''})^{(7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7)}(T_{37},t)$ are seventh augmentation coefficients for category 1, 2 and 3<br>$+(a_{40}^{''})^{(8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8)}(T_{41},t)$ are eight augmentation coefficients | |

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| for category 1, 2 and 3<br>$+(a_{44}^{''})^{(9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9)}(T_{45},t)$, $+(a_{46}^{''})^{(9,9,9)}(T_{45},t)$ are ninth augmentation coefficients for category 1, 2 and 3 | |
| $\frac{dT_{20}}{dt} = (b_{20})^{(3)}T_{21} - \begin{bmatrix} (b_{20}^{'})^{(3)}\,-(b_{20}^{''})^{(3)}(G_{23},t) & -(b_{16}^{''})^{(2,2,2)}(G_{19},t) & -(b_{13}^{''})^{(1,1,1,)}(G,t) \\ -(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t) & -(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t) & -(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t) \\ -(b_{36}^{''})^{(7,7,7,7)}(G_{39},t) & -(b_{40}^{''})^{(8,8,8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9,9)}(G_{47},t) \end{bmatrix} T_{20}$ | *70* |
| $\frac{dT_{21}}{dt} = (b_{21})^{(3)}T_{20} - \begin{bmatrix} (b_{21}^{'})^{(3)}\,-(b_{21}^{''})^{(3)}(G_{23},t) & -(b_{17}^{''})^{(2,2,2)}(G_{19},t) & -(b_{14}^{''})^{(1,1,1,)}(G,t) \\ -(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t) & -(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t) & -(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t) \\ -(b_{37}^{''})^{(7,7,7,7)}(G_{39},t) & -(b_{41}^{''})^{(8,8,8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9,9)}(G_{47},t) \end{bmatrix} T_{21}$ | *71* |
| $\frac{dT_{22}}{dt} = (b_{22})^{(3)}T_{21} - \begin{bmatrix} (b_{22}^{'})^{(3)}\,-(b_{22}^{''})^{(3)}(G_{23},t) & -(b_{18}^{''})^{(2,2,2)}(G_{19},t) & -(b_{15}^{''})^{(1,1,1,)}(G,t) \\ -(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t) & -(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t) & -(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t) \\ -(b_{38}^{''})^{(7,7,7,7)}(G_{39},t) & -(b_{42}^{''})^{(8,8,8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9)}(G_{47},t) \end{bmatrix} T_{22}$ | *72* |
| $-(b_{20}^{''})^{(3)}(G_{23},t)$, $-(b_{21}^{''})^{(3)}(G_{23},t)$, $-(b_{22}^{''})^{(3)}(G_{23},t)$ *are first detrition coefficients for category 1, 2 and 3*<br>$-(b_{16}^{''})^{(2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2)}(G_{19},t)$ *are second detrition coefficients for category 1, 2 and 3*<br>$-(b_{13}^{''})^{(1,1,1,)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,)}(G,t)$, $-(b_{15}^{''})^{(1,1,1,)}(G,t)$ *are third detrition coefficients for category 1,2 and 3*<br>$-(b_{24}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4)}(G_{27},t)$ *are fourth detrition coefficients for category 1, 2 and 3*<br>$-(b_{28}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5)}(G_{31},t)$ *are fifth detrition coefficients for category 1, 2 and 3*<br>$-(b_{32}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6)}(G_{35},t)$ *are sixth detrition coefficients for category 1, 2 and 3*<br>$-(b_{36}^{''})^{(7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7)}(G_{39},t)$ $-(b_{38}^{''})^{(7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3<br>$-(b_{40}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8)}(G_{43},t)$ are eight detrition coefficients for category 1, 2 and 3<br>$-(b_{46}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3 | |
| $\frac{dG_{24}}{dt} = (a_{24})^{(4)}G_{25} - \begin{bmatrix} (a_{24}^{'})^{(4)}\,+(a_{24}^{''})^{(4)}(T_{25},t) & +(a_{28}^{''})^{(5,5,)}(T_{29},t) & +(a_{32}^{''})^{(6,6,)}(T_{33},t) \\ +(a_{13}^{''})^{(1,1,1,1)}(T_{14},t) & +(a_{16}^{''})^{(2,2,2,2)}(T_{17},t) & +(a_{20}^{''})^{(3,3,3,3)}(T_{21},t) \\ +(a_{36}^{''})^{(7,7,7,7,7)}(T_{37},t) & +(a_{40}^{''})^{(8,8,8,8,8)}(T_{41},t) & +(a_{44}^{''})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{24}$ | *73* |

$$\frac{dG_{25}}{dt} = (a_{25})^{(4)}G_{24} - \begin{bmatrix} (a_{25}^{'})^{(4)} & +(a_{25}^{''})^{(4)}(T_{25},t) & +(a_{29}^{''})^{(5,5,)}(T_{29},t) & +(a_{33}^{''})^{(6,6)}(T_{33},t) \\ & +(a_{14}^{''})^{(1,1,1,1)}(T_{14},t) & +(a_{17}^{''})^{(2,2,2,2)}(T_{17},t) & +(a_{21}^{''})^{(3,3,3,3)}(T_{21},t) \\ & +(a_{37}^{''})^{(7,7,7,7,7)}(T_{37},t) & +(a_{41}^{''})^{(8,8,8,8,8)}(T_{41},t) & +(a_{45}^{''})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{25} \quad (74)$$

$$\frac{dG_{26}}{dt} = (a_{26})^{(4)}G_{25} - \begin{bmatrix} (a_{26}^{'})^{(4)} & +(a_{26}^{''})^{(4)}(T_{25},t) & +(a_{30}^{''})^{(5,5,)}(T_{29},t) & +(a_{34}^{''})^{(6,6,)}(T_{33},t) \\ & +(a_{15}^{''})^{(1,1,1,1)}(T_{14},t) & +(a_{18}^{''})^{(2,2,2,2)}(T_{17},t) & +(a_{22}^{''})^{(3,3,3,3)}(T_{21},t) \\ & +(a_{38}^{''})^{(7,7,7,7,7)}(T_{37},t) & +(a_{42}^{''})^{(8,8,8,8,8)}(T_{41},t) & +(a_{46}^{''})^{(9,9,9,9)}(T_{45},t) \end{bmatrix} G_{26} \quad (75)$$

$(a_{24}^{''})^{(4)}(T_{25},t)$, $(a_{25}^{''})^{(4)}(T_{25},t)$, $(a_{26}^{''})^{(4)}(T_{25},t)$ *are first augmentation coefficients category* 1, 2 3

$+(a_{28}^{''})^{(5,5,)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,)}(T_{29},t)$ *are second augmentation coefficient for category* 1, 2 *and* 3

$+(a_{32}^{''})^{(6,6,)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,)}(T_{33},t)$ *are third augmentation coefficient for category* 1, 2 *and* 3

$+(a_{13}^{''})^{(1,1,1,1)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1)}(T_{14},t)$ *are fourth augmentation coefficients for category* 1, 2 *and* 3

$+(a_{16}^{''})^{(2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2)}(T_{17},t)$ *are fifth augmentation coefficients for category* 1, 2 *and* 3

$+(a_{20}^{''})^{(3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3)}(T_{21},t)$ *are sixth augmentation coefficients for category* 1, 2 *and* 3

$+(a_{36}^{''})^{(7,7,7,7,7)}(T_{37},t)$, $+(a_{37}^{''})^{(7,7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7,7)}(T_{37},t)$ *are seventh augmentation coefficients for category* 1, 2 *and* 3

$+(a_{40}^{''})^{(8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8,8)}(T_{41},t)$ *are eighth augmentation coefficients for category* 1, 2 *and* 3

$+(a_{46}^{''})^{(9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9)}(T_{45},t)$ are ninth detrition coefficients for category 1 2 3

$$\frac{dT_{24}}{dt} = (b_{24})^{(4)}T_{25} - \begin{bmatrix} (b_{24}^{'})^{(4)} & -(b_{24}^{''})^{(4)}(G_{27},t) & -(b_{28}^{''})^{(5,5,)}(G_{31},t) & -(b_{32}^{''})^{(6,6,)}(G_{35},t) \\ & -(b_{13}^{''})^{(1,1,1,1)}(G,t) & -(b_{16}^{''})^{(2,2,2,2)}(G_{19},t) & -(b_{20}^{''})^{(3,3,3,3)}(G_{23},t) \\ & -(b_{36}^{''})^{(7,7,7,7,7)}(G_{39},t) & -(b_{40}^{''})^{(8,8,8,8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{24} \quad (76)$$

$$\frac{dT_{25}}{dt} = (b_{25})^{(4)}T_{24} - \begin{bmatrix} (b_{25}^{'})^{(4)} & -(b_{25}^{''})^{(4)}(G_{27},t) & -(b_{29}^{''})^{(5,5,)}(G_{31},t) & -(b_{33}^{''})^{(6,6,)}(G_{35},t) \\ & -(b_{14}^{''})^{(1,1,1,1)}(G,t) & -(b_{17}^{''})^{(2,2,2,2)}(G_{19},t) & -(b_{21}^{''})^{(3,3,3,3)}(G_{23},t) \\ & -(b_{37}^{''})^{(7,7,7,7,7)}(G_{39},t) & -(b_{41}^{''})^{(8,8,8,8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{25} \quad (77)$$

$$\frac{dT_{26}}{dt} = (b_{26})^{(4)}T_{25} - \begin{bmatrix} (b_{26}^{'})^{(4)} & -(b_{26}^{''})^{(4)}(G_{27},t) & -(b_{30}^{''})^{(5,5,)}(G_{31},t) & -(b_{34}^{''})^{(6,6,)}(G_{35},t) \\ & -(b_{15}^{''})^{(1,1,1,1)}(G,t) & -(b_{18}^{''})^{(2,2,2,2)}(G_{19},t) & -(b_{22}^{''})^{(3,3,3,3)}(G_{23},t) \\ & -(b_{38}^{''})^{(7,7,7,7,7)}(G_{39},t) & -(b_{42}^{''})^{(8,8,8,8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9,9)}(G_{47},t) \end{bmatrix} T_{26} \quad (78)$$

*Where* $-(b_{24}^{''})^{(4)}(G_{27},t)$, $-(b_{25}^{''})^{(4)}(G_{27},t)$, $-(b_{26}^{''})^{(4)}(G_{27},t)$ *are first detrition coefficients for category* 1, 2 *and* 3

$-(b_{28}^{''})^{(5,5,)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,)}(G_{31},t)$ *are second detrition coefficients for category* 1, 2 *and* 3

$-(b_{32}^{''})^{(6,6,)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,)}(G_{35},t)$ *are third detrition coefficients for category* 1, 2 *and* 3

$-(b_{13}^{''})^{(1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1)}(G,t)$, $-(b_{15}^{''})^{(1,1,1,1)}(G,t)$ *are fourth detrition coefficients for category* 1, 2 *and* 3

$-(b_{16}^{''})^{(2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2)}(G_{19},t)$ *are fifth detrition coefficients for category* 1, 2 *and* 3

$-(b_{20}^{''})^{(3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3)}(G_{23},t)$ *are sixth detrition coefficients for category* 1, 2 *and* 3

$-(b_{36}^{''})^{(7,7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7,7)}(G_{39},t)$ *are seventh detrition coefficients for category* 1, 2 *and* 3

$-(b_{40}^{''})^{(8,8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8,8)}(G_{43},t)$ *are eighth detrition coefficients for category* 1, 2 *and* 3

$-(b_{46}^{''})^{(9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1 2 3

$$\frac{dG_{28}}{dt} = (a_{28})^{(5)}G_{29} - \left[\begin{array}{ccc} (a_{28}^{'})^{(5)} & +(a_{28}^{''})^{(5)}(T_{29},t) & +(a_{24}^{''})^{(4,4,)}(T_{25},t) \quad +(a_{32}^{''})^{(6,6,6)}(T_{33},t) \\ +(a_{13}^{''})^{(1,1,1,1,1)}(T_{14},t) & +(a_{16}^{''})^{(2,2,2,2,2)}(T_{17},t) & +(a_{20}^{''})^{(3,3,3,3,3)}(T_{21},t) \\ +(a_{36}^{''})^{(7,7,7,7,7)}(T_{37},t) & +(a_{40}^{''})^{(8,8,8,8,8)}(T_{41},t) & +(a_{44}^{''})^{(9,9,9,9,9)}(T_{45},t) \end{array}\right] G_{28} \quad (79)$$

$$\frac{dG_{29}}{dt} = (a_{29})^{(5)}G_{28} - \left[\begin{array}{ccc} (a_{29}^{'})^{(5)} & +(a_{29}^{''})^{(5)}(T_{29},t) & +(a_{25}^{''})^{(4,4,)}(T_{25},t) \quad +(a_{33}^{''})^{(6,6,6)}(T_{33},t) \\ +(a_{14}^{''})^{(1,1,1,1,1)}(T_{14},t) & +(a_{17}^{''})^{(2,2,2,2,2)}(T_{17},t) & +(a_{21}^{''})^{(3,3,3,3,3)}(T_{21},t) \\ +(a_{37}^{''})^{(7,7,7,7,7)}(T_{37},t) & +(a_{41}^{''})^{(8,8,8,8,8)}(T_{41},t) & +(a_{45}^{''})^{(9,9,9,9,9)}(T_{45},t) \end{array}\right] G_{29} \quad (80)$$

$$\frac{dG_{30}}{dt} = (a_{30})^{(5)}G_{29} - \left[\begin{array}{ccc} (a_{30}^{'})^{(5)} & +(a_{30}^{''})^{(5)}(T_{29},t) & +(a_{26}^{''})^{(4,4,)}(T_{25},t) \quad +(a_{34}^{''})^{(6,6,6)}(T_{33},t) \\ +(a_{15}^{''})^{(1,1,1,1,1)}(T_{14},t) & +(a_{18}^{''})^{(2,2,2,2,2)}(T_{17},t) & +(a_{22}^{''})^{(3,3,3,3,3)}(T_{21},t) \\ +(a_{38}^{''})^{(7,7,7,7,7)}(T_{37},t) & +(a_{42}^{''})^{(8,8,8,8,8)}(T_{41},t) & +(a_{46}^{''})^{(9,9,9,9,9)}(T_{45},t) \end{array}\right] G_{30} \quad (81)$$

*Where* $+(a_{28}^{''})^{(5)}(T_{29},t)$, $+(a_{29}^{''})^{(5)}(T_{29},t)$, $+(a_{30}^{''})^{(5)}(T_{29},t)$ *are first augmentation coefficients for category* 1, 2 *and* 3

*And* $+(a_{24}^{''})^{(4,4,)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,)}(T_{25},t)$ *are second augmentation coefficient for category* 1, 2 *and* 3

$+(a_{32}^{''})^{(6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6)}(T_{33},t)$ *are third augmentation coefficient for category* 1, 2 *and* 3

$+(a_{13}^{''})^{(1,1,1,1,1)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1,1)}(T_{14},t)$ are fourth augmentation coefficients for category 1,2, and 3

$+(a_{16}^{''})^{(2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2)}(T_{17},t)$ are fifth augmentation coefficients for category 1,2,and 3

$+(a_{20}^{''})^{(3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3)}(T_{21},t)$ are sixth augmentation

| | |
|---|---|
| coefficients for category 1,2, 3<br>$+(a_{36}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)$ , $+(a_{37}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)$, $+(a_{38}^{''})^{(7,7,7,7,7,7,7)}(T_{37},t)$ are seventh augmentation coefficients for category 1,2, 3<br>$+(a_{40}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8,8,8,8)}(T_{41},t)$ are eighth augmentation coefficients for category 1,2, 3<br>$+(a_{46}^{''})^{(9,9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficients for category 1,2, 3 | |
| $\frac{dT_{28}}{dt} = (b_{28})^{(5)}T_{29} - \begin{bmatrix} (b_{28}^{'})^{(5)} & -(b_{28}^{''})^{(5)}(G_{31},t) & -(b_{24}^{''})^{(4,4,)}(G_{27},t) & -(b_{32}^{''})^{(6,6,6)}(G_{35},t) \\ & -(b_{13}^{''})^{(1,1,1,1,1)}(G,t) & -(b_{16}^{''})^{(2,2,2,2,2)}(G_{19},t) & -(b_{20}^{''})^{(3,3,3,3,3)}(G_{23},t) \\ & -(b_{36}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b_{40}^{''})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{28}$ | *82* |
| $\frac{dT_{29}}{dt} = (b_{29})^{(5)}T_{28} - \begin{bmatrix} (b_{29}^{'})^{(5)} & -(b_{29}^{''})^{(5)}(G_{31},t) & -(b_{25}^{''})^{(4,4,)}(G_{27},t) & -(b_{33}^{''})^{(6,6,6)}(G_{35},t) \\ & -(b_{14}^{''})^{(1,1,1,1,1)}(G,t) & -(b_{17}^{''})^{(2,2,2,2,2)}(G_{19},t) & -(b_{21}^{''})^{(3,3,3,3,3)}(G_{23},t) \\ & -(b_{37}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b_{41}^{''})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{29}$ | *83* |
| $\frac{dT_{30}}{dt} = (b_{30})^{(5)}T_{29} - \begin{bmatrix} (b_{30}^{'})^{(5)} & -(b_{30}^{''})^{(5)}(G_{31},t) & -(b_{26}^{''})^{(4,4,)}(G_{27},t) & -(b_{34}^{''})^{(6,6,6)}(G_{35},t) \\ & -(b_{15}^{''})^{(1,1,1,1,1,)}(G,t) & -(b_{18}^{''})^{(2,2,2,2,2)}(G_{19},t) & -(b_{22}^{''})^{(3,3,3,3,3)}(G_{23},t) \\ & -(b_{38}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b_{42}^{''})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{30}$ | *84* |
| *where* $-(b_{28}^{''})^{(5)}(G_{31},t)$ , $-(b_{29}^{''})^{(5)}(G_{31},t)$ , $-(b_{30}^{''})^{(5)}(G_{31},t)$ *are first detrition coefficients for category* 1, 2 *and* 3<br>$-(b_{24}^{''})^{(4,4,)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,)}(G_{27},t)$ *are second detrition coefficients for category* 1,2 *and* 3<br>$-(b_{32}^{''})^{(6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6)}(G_{35},t)$ *are third detrition coefficients for category* 1,2 *and* 3<br>$-(b_{13}^{''})^{(1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1)}(G,t)$ , $-(b_{15}^{''})^{(1,1,1,1,1,)}(G,t)$ are fourth detrition coefficients for category 1,2, and 3<br>$-(b_{16}^{''})^{(2,2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2,2)}(G_{19},t)$ are fifth detrition coefficients for category 1,2, and 3<br>$-(b_{20}^{''})^{(3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3)}(G_{23},t)$ are sixth detrition coefficients for category 1,2, and 3<br>$-(b_{36}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7,7,7,7)}(G_{39},t)$ are seventh detrition coefficients for category 1,2, and 3<br>$-(b_{42}^{''})^{(8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{40}^{''})^{(8,8,8,8,8,8,8)}(G_{43},t)$ are eighth detrition coefficients for category 1,2, and 3<br>$-(b_{46}^{''})^{(9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1,2, and 3 | |

$$\frac{dG_{32}}{dt} = (a_{32})^{(6)} G_{33} - \begin{bmatrix} (a'_{32})^{(6)} & +(a''_{32})^{(6)}(T_{33},t) & +(a''_{28})^{(5,5,5)}(T_{29},t) & +(a''_{24})^{(4,4,4,)}(T_{25},t) \\ & +(a''_{13})^{(1,1,1,1,1,1)}(T_{14},t) & +(a''_{16})^{(2,2,2,2,2,2)}(T_{17},t) & +(a''_{20})^{(3,3,3,3,3,3)}(T_{21},t) \\ & +(a''_{36})^{(7,7,7,7,7,7,7)}(T_{37},t) & +(a''_{40})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a''_{44})^{(9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{32} \quad (85)$$

$$\frac{dG_{33}}{dt} = (a_{33})^{(6)} G_{32} - \begin{bmatrix} (a'_{33})^{(6)} & +(a''_{33})^{(6)}(T_{33},t) & +(a''_{29})^{(5,5,5)}(T_{29},t) & +(a''_{25})^{(4,4,4,)}(T_{25},t) \\ & +(a''_{14})^{(1,1,1,1,1,1)}(T_{14},t) & +(a''_{17})^{(2,2,2,2,2,2)}(T_{17},t) & +(a''_{21})^{(3,3,3,3,3,3)}(T_{21},t) \\ & +(a''_{37})^{(7,7,7,7,7,7,7)}(T_{37},t) & +(a''_{41})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a''_{45})^{(9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{33} \quad (86)$$

$$\frac{dG_{34}}{dt} = (a_{34})^{(6)} G_{33} - \begin{bmatrix} (a'_{34})^{(6)} & +(a''_{34})^{(6)}(T_{33},t) & +(a''_{30})^{(5,5,5)}(T_{29},t) & +(a''_{26})^{(4,4,4,)}(T_{25},t) \\ & +(a''_{15})^{(1,1,1,1,1,1)}(T_{14},t) & +(a''_{18})^{(2,2,2,2,2,2)}(T_{17},t) & +(a''_{22})^{(3,3,3,3,3,3)}(T_{21},t) \\ & +(a''_{38})^{(7,7,7,7,7,7,7)}(T_{37},t) & +(a''_{42})^{(8,8,8,8,8,8,8)}(T_{41},t) & +(a''_{46})^{(9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{34} \quad (87)$$

$+(a''_{32})^{(6)}(T_{33},t)$, $+(a''_{33})^{(6)}(T_{33},t)$, $+(a''_{34})^{(6)}(T_{33},t)$ *are first augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{28})^{(5,5,5)}(T_{29},t)$, $+(a''_{29})^{(5,5,5)}(T_{29},t)$, $+(a''_{30})^{(5,5,5)}(T_{29},t)$ *are second augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{24})^{(4,4,4,)}(T_{25},t)$, $+(a''_{25})^{(4,4,4,)}(T_{25},t)$, $+(a''_{26})^{(4,4,4,)}(T_{25},t)$ *are third augmentation coefficients for category* 1, 2 *and* 3

$+(a''_{13})^{(1,1,1,1,1,1)}(T_{14},t)$, $+(a''_{14})^{(1,1,1,1,1,1)}(T_{14},t)$, $+(a''_{15})^{(1,1,1,1,1,1)}(T_{14},t)$ - are fourth augmentation coefficients

$+(a''_{16})^{(2,2,2,2,2,2)}(T_{17},t)$, $+(a''_{17})^{(2,2,2,2,2,2)}(T_{17},t)$, $+(a''_{18})^{(2,2,2,2,2,2)}(T_{17},t)$ - fifth augmentation coefficients

$+(a''_{20})^{(3,3,3,3,3,3)}(T_{21},t)$, $+(a''_{21})^{(3,3,3,3,3,3)}(T_{21},t)$, $+(a''_{22})^{(3,3,3,3,3,3)}(T_{21},t)$ sixth augmentation coefficients

$+(a''_{36})^{(7,7,7,7,7,7,7)}(T_{37},t)$, $+(a''_{37})^{(7,7,7,7,7,7,7)}(T_{37},t)$, $+(a''_{38})^{(7,7,7,7,7,7,7)}(T_{37},t)$ seventh augmentation coefficients

$+(a''_{40})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a''_{41})^{(8,8,8,8,8,8,8)}(T_{41},t)$, $+(a''_{42})^{(8,8,8,8,8,8,8)}(T_{41},t)$ Eighth augmentation coefficients

$+(a''_{44})^{(9,9,9,9,9,9)}(T_{45},t)$, $+(a''_{45})^{(9,9,9,9,9,9)}(T_{45},t)$, $+(a''_{46})^{(9,9,9,9,9,9)}(T_{45},t)$ ninth augmentation coefficients

$$\frac{dT_{32}}{dt} = (b_{32})^{(6)} T_{33} - \begin{bmatrix} (b'_{32})^{(6)} & -(b''_{32})^{(6)}(G_{35},t) & -(b''_{28})^{(5,5,5)}(G_{31},t) & -(b''_{24})^{(4,4,4,)}(G_{27},t) \\ & -(b''_{13})^{(1,1,1,1,1,1)}(G,t) & -(b''_{16})^{(2,2,2,2,2,2)}(G_{19},t) & -(b''_{20})^{(3,3,3,3,3,3)}(G_{23},t) \\ & -(b''_{36})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b''_{40})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b''_{44})^{(9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{32} \quad (88)$$

$$\frac{dT_{33}}{dt} = (b_{33})^{(6)} T_{32} - \begin{bmatrix} (b'_{33})^{(6)} & -(b''_{33})^{(6)}(G_{35},t) & -(b''_{29})^{(5,5,5)}(G_{31},t) & -(b''_{25})^{(4,4,4,)}(G_{27},t) \\ & -(b''_{14})^{(1,1,1,1,1,1)}(G,t) & -(b''_{17})^{(2,2,2,2,2,2)}(G_{19},t) & -(b''_{21})^{(3,3,3,3,3,3)}(G_{23},t) \\ & -(b''_{37})^{(7,7,7,7,7,7,7)}(G_{39},t) & -(b''_{41})^{(8,8,8,8,8,8,8)}(G_{43},t) & -(b''_{45})^{(9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{33} \quad (89)$$

$$\frac{dT_{34}}{dt} = (b_{34})^{(6)}T_{33} - \begin{bmatrix} (b'_{34})^{(6)}\boxed{-(b''_{34})^{(6)}(G_{35},t)} & \boxed{-(b''_{30})^{(5,5,5)}(G_{31},t)} & \boxed{-(b''_{26})^{(4,4,4,)}(G_{27},t)} \\ \boxed{-(b''_{15})^{(1,1,1,1,1,1)}(G,t)} & \boxed{-(b''_{18})^{(2,2,2,2,2,2)}(G_{19},t)} & \boxed{-(b''_{22})^{(3,3,3,3,3,3)}(G_{23},t)} \\ \boxed{-(b''_{38})^{(7,7,7,7,7,7,7)}(G_{39},t)} & \boxed{-(b''_{42})^{(8,8,8,8,8,8,8)}(G_{43},t)} & \boxed{-(b''_{46})^{(9,9,9,9,9,9)}(G_{47},t)} \end{bmatrix} T_{34} \quad (90)$$

$\boxed{-(b''_{32})^{(6)}(G_{35},t)}, \boxed{-(b''_{33})^{(6)}(G_{35},t)}, \boxed{-(b''_{34})^{(6)}(G_{35},t)}$ *are first detrition coefficients for category* 1, 2 *and* 3

$\boxed{-(b''_{28})^{(5,5,5)}(G_{31},t)}, \boxed{-(b''_{29})^{(5,5,5)}(G_{31},t)}, \boxed{-(b''_{30})^{(5,5,5)}(G_{31},t)}$ *are second detrition coefficients for category* 1, 2 *and* 3

$\boxed{-(b''_{24})^{(4,4,4,)}(G_{27},t)}, \boxed{-(b''_{25})^{(4,4,4,)}(G_{27},t)}, \boxed{-(b''_{26})^{(4,4,4,)}(G_{27},t)}$ *are third detrition coefficients for category* 1,2 *and* 3

$\boxed{-(b''_{13})^{(1,1,1,1,1,1)}(G,t)}, \boxed{-(b''_{14})^{(1,1,1,1,1,1)}(G,t)}, \boxed{-(b''_{15})^{(1,1,1,1,1,1)}(G,t)}$ are fourth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{16})^{(2,2,2,2,2,2)}(G_{19},t)}, \boxed{-(b''_{17})^{(2,2,2,2,2,2)}(G_{19},t)}, \boxed{-(b''_{18})^{(2,2,2,2,2,2)}(G_{19},t)}$ are fifth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{20})^{(3,3,3,3,3,3)}(G_{23},t)}, \boxed{-(b''_{21})^{(3,3,3,3,3,3)}(G_{23},t)}, \boxed{-(b''_{22})^{(3,3,3,3,3,3)}(G_{23},t)}$ are sixth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{36})^{(7,7,7,7,7,7,7)}(G_{39},t)}, \boxed{-(b''_{37})^{(7,7,7,7,7,7,7)}(G_{39},t)}, \boxed{-(b''_{38})^{(7,7,7,7,7,7,7)}(G_{39},t)}$ are seventh detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{40})^{(8,8,8,8,8,8,8,8)}(G_{43},t)}, \boxed{-(b''_{41})^{(8,8,8,8,8,8,8,8)}(G_{43},t)}, \boxed{-(b''_{42})^{(8,8,8,8,8,8,8,8)}(G_{43},t)}$ are eighth detrition coefficients for category 1, 2, and 3

$\boxed{-(b''_{46})^{(9,9,9,9,9,9)}(G_{47},t)}, \boxed{-(b''_{45})^{(9,9,9,9,9,9)}(G_{47},t)}, \boxed{-(b''_{44})^{(9,9,9,9,9,9)}(G_{47},t)}$ are ninth detrition coefficients for category 1, 2, and 3

$$\frac{dG_{36}}{dt} = (a_{36})^{(7)}G_{37} - \begin{bmatrix} (a'_{36})^{(7)}\boxed{+(a''_{36})^{(7)}(T_{37},t)} & \boxed{+(a''_{16})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{20})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{24})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{28})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{32})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{40})^{(8,8,8,8,8,8,8,8,)}(T_{41},t)} & \boxed{+(a''_{44})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{13} \quad (91)$$

$$\frac{dG_{37}}{dt} = (a_{37})^{(7)}G_{36} - \begin{bmatrix} (a'_{37})^{(7)}\boxed{+(a''_{37})^{(7)}(T_{37},t)} & \boxed{+(a''_{17})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{21})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{25})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{29})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{33})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{13})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{41})^{(8,8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{45})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{14} \quad (92)$$

$$\frac{dG_{38}}{dt} = (a_{38})^{(7)}G_{37} - \begin{bmatrix} (a'_{38})^{(7)}\boxed{+(a''_{38})^{(7)}(T_{37},t)} & \boxed{+(a''_{18})^{(2,2,2,2,2,2,2)}(T_{17},t)} & \boxed{+(a''_{22})^{(3,3,3,3,3,3,3)}(T_{21},t)} \\ \boxed{+(a''_{26})^{(4,4,4,4,4,4,4)}(T_{25},t)} & \boxed{+(a''_{30})^{(5,5,5,5,5,5,5)}(T_{29},t)} & \boxed{+(a''_{34})^{(6,6,6,6,6,6,6)}(T_{33},t)} \\ \boxed{+(a''_{15})^{(1,1,1,1,1,1,1)}(T_{14},t)} & \boxed{+(a''_{42})^{(8,8,8,8,8,8,8,8)}(T_{41},t)} & \boxed{+(a''_{46})^{(9,9,9,9,9,9,9)}(T_{45},t)} \end{bmatrix} G_{15} \quad (93)$$

Where $(a_{36}^{''})^{(7)}(T_{37},t)$, $(a_{37}^{''})^{(7)}(T_{37},t)$, $(a_{38}^{''})^{(7)}(T_{37},t)$ are first augmentation coefficients for category 1, 2 and 3

$+(a_{16}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2,2,2)}(T_{17},t)$ are second augmentation coefficient for category 1, 2 and 3

$+(a_{20}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3,3,3)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3

$+(a_{24}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3

$+(a_{28}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficient for category 1, 2 and 3

$+(a_{32}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation coefficient for category 1, 2 and 3

$+(a_{13}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{13}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1,1,1,1)}(T_{14},t)$ are seventh augmentation coefficient for category 1, 2 and 3

$+(a_{42}^{''})^{(8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{40}^{''})^{(8,8,8,8,8,8,8,8,)}(T_{41},t)$ are eighth augmentation coefficient for 1,2,3

$+(a_{46}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{45}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t)$, $+(a_{44}^{''})^{(9,9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{36}}{dt} = (b_{36})^{(7)}T_{37} - \begin{bmatrix} (b_{36}')^{(7)}-(b_{36}^{''})^{(7)}(G_{39},t) & -(b_{16}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t) & -(b_{20}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t) \\ -(b_{24}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t) & -(b_{28}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t) & -(b_{32}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t) \\ -(b_{13}^{''})^{(1,1,1,1,1,1,1)}(G,t) & -(b_{40}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t) & -(b_{44}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{13} \quad (94)$$

$$\frac{dT_{37}}{dt} = (b_{37})^{(7)}T_{36} - \begin{bmatrix} (b_{37}')^{(7)}-(b_{37}^{''})^{(7)}(G_{39},t) & -(b_{17}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t) & -(b_{21}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t) \\ -(b_{25}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t) & -(b_{29}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t) & -(b_{33}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t) \\ -(b_{14}^{''})^{(1,1,1,1,1,1,1)}(G,t) & -(b_{41}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t) & -(b_{45}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{14}$$

$$\frac{dT_{38}}{dt} = (b_{38})^{(7)}T_{37} - \begin{bmatrix} (b_{38}')^{(7)}-(b_{38}^{''})^{(7)}(G_{39},t) & -(b_{18}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t) & -(b_{22}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t) \\ -(b_{26}^{''})^{(4,4,4,4,4,4,4)}(G_{27},t) & -(b_{30}^{''})^{(5,5,5,5,5,5,5)}(G_{31},t) & -(b_{34}^{''})^{(6,6,6,6,6,6,6)}(G_{35},t) \\ -(b_{15}^{''})^{(1,1,1,1,1,1,1)}(G,t) & -(b_{42}^{''})^{(8,8,8,8,8,8,8,8)}(G_{43},t) & -(b_{46}^{''})^{(9,9,9,9,9,9,9)}(G_{47},t) \end{bmatrix} T_{15}$$

Where $-(b_{36}^{''})^{(7)}(G_{39},t)$, $-(b_{37}^{''})^{(7)}(G_{39},t)$, $-(b_{38}^{''})^{(7)}(G_{39},t)$ are first detrition coefficients for category 1, 2 and 3

$-(b_{16}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2,2,2,2)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3

| | |
|---|---|
| $-(b_{20}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3,3,3)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3<br>$-(b_{24}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3<br>$-(b_{28}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3<br>$-(b_{32}^{''})^{(6,6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6,6,6)}(G_{35},t)$ are sixth detrition coefficients for category 1, 2 and 3<br>$-(b_{15}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{13}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$ are seventh detrition coefficients for category 1, 2 and 3<br>$-(b_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)$, $-(b_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)$ are eighth detrition coefficients for category 1, 2 and 3<br>$-(b_{46}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{44}^{''})^{(9,9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3 | |
| $\frac{dG_{40}}{dt} = (a_{40})^{(8)}G_{41} - \begin{bmatrix} (a_{40}^{'})^{(8)}+(a_{40}^{''})^{(8)}(T_{41},t) & +(a_{16}^{''})^{(2,2,2,2,2,2,2,2)}(T_{17},t) & +(a_{20}^{''})^{(3,3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{24}^{''})^{(4,4,4,4,4,4,4,4)}(T_{25},t) & +(a_{28}^{''})^{(5,5,5,5,5,5,5,5)}(T_{29},t) & +(a_{32}^{''})^{(6,6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{13}^{''})^{(1,1,1,1,1,1,1,1)}(T_{14},t) & +(a_{36}^{''})^{(7,7,7,7,7,7,7,7)}(T_{37},t) & +(a_{44}^{''})^{(9,9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{13}$ | *95* |
| $\frac{dG_{41}}{dt} = (a_{41})^{(8)}G_{40} - \begin{bmatrix} (a_{41}^{'})^{(8)}+(a_{41}^{''})^{(8)}(T_{41},t) & +(a_{17}^{''})^{(2,2,2,2,2,2,2,2)}(T_{17},t) & +(a_{21}^{''})^{(3,3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{25}^{''})^{(4,4,4,4,4,4,4,4)}(T_{25},t) & +(a_{29}^{''})^{(5,5,5,5,5,5,5,5)}(T_{29},t) & +(a_{33}^{''})^{(6,6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{13}^{''})^{(1,1,1,1,1,1,1,1)}(T_{14},t) & +(a_{37}^{''})^{(7,7,7,7,7,7,7,7)}(T_{37},t) & +(a_{45}^{''})^{(9,9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{14}$ | |
| $\frac{dG_{42}}{dt} = (a_{42})^{(8)}G_{41} - \begin{bmatrix} (a_{42}^{'})^{(8)}+(a_{42}^{''})^{(8)}(T_{41},t) & +(a_{18}^{''})^{(2,2,2,2,2,2,2,2)}(T_{17},t) & +(a_{22}^{''})^{(3,3,3,3,3,3,3,3)}(T_{21},t) \\ +(a_{26}^{''})^{(4,4,4,4,4,4,4,4)}(T_{25},t) & +(a_{30}^{''})^{(5,5,5,5,5,5,5,5)}(T_{29},t) & +(a_{34}^{''})^{(6,6,6,6,6,6,6,6)}(T_{33},t) \\ +(a_{15}^{''})^{(1,1,1,1,1,1,1,1)}(T_{14},t) & +(a_{38}^{''})^{(7,7,7,7,7,7,7,7)}(T_{37},t) & +(a_{46}^{''})^{(9,9,9,9,9,9,9,9)}(T_{45},t) \end{bmatrix} G_{15}$ | |
| Where $+(a_{40}^{''})^{(8)}(T_{41},t)$, $+(a_{41}^{''})^{(8)}(T_{41},t)$, $+(a_{42}^{''})^{(8)}(T_{41},t)$ are first augmentation coefficients for category 1, 2 and 3<br>$+(a_{16}^{''})^{(2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2,2,2,2)}(T_{17},t)$ are second augmentation coefficient for category 1, 2 and 3<br>$+(a_{20}^{''})^{(3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3,3,3,3)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3 | |

$+(a_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$ , $+(a_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$ , $+(a_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3

$+(a_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$ , $+(a_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficient for category 1, 2 and 3

$+(a_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$ , $+(a_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$ are sixth augmentation coefficient for category 1, 2 and 3

$+(a_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)$ $+(a_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)$ $+(a_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)$ are seventh augmentation coefficient for 1,2,3

$+(a_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)$ ,$+(a_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)$ ,$+(a_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)$ are eighth augmentation coefficient for 1,2,3

$+(a_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ ,$+(a_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ ,$+(a_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(T_{45},t)$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{40}}{dt} =$$

$$(b_{40})^{(8)}T_{41} - \begin{bmatrix} (b_{40}^{'})^{(8)}\left[-(b_{40}^{''})^{(8)}(G_{43},t)\right] & \left[-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)\right] & \left[-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)\right] \\ \left[-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)\right] & \left[-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)\right] & \left[-(b_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)\right] \\ \left[-(b_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)\right] & \left[-(b_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)\right] & \left[-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{13}$$

$$\frac{dT_{41}}{dt} =$$

$$(b_{41})^{(8)}T_{40} - \begin{bmatrix} (b_{41}^{'})^{(8)}\left[-(b_{41}^{''})^{(8)}(G_{43},t)\right] & \left[-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)\right] & \left[-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)\right] \\ \left[-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)\right] & \left[-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)\right] & \left[-(b_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)\right] \\ \left[-(b_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)\right] & \left[-(b_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)\right] & \left[-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{14}$$

$$\frac{dT_{42}}{dt} =$$

$$(b_{42})^{(8)}T_{41} - \begin{bmatrix} (b_{42}^{'})^{(8)}\left[-(b_{42}^{''})^{(8)}(G_{43},t)\right] & \left[-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)\right] & \left[-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)\right] \\ \left[-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)\right] & \left[-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)\right] & \left[-(b_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(G_{35},t)\right] \\ \left[-(b_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(G,t)\right] & \left[-(b_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(G_{39},t)\right] & \left[-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)\right] \end{bmatrix} T_{15}$$

Where $-(b_{36}^{''})^{(7)}(G_{39},t)$ , $-(b_{37}^{''})^{(7)}(G_{39},t)$ , $-(b_{38}^{''})^{(7)}(G_{39},t)$ are first detrition coefficients for category 1, 2 and 3

$-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)$ , $-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)$ , $-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3

$-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)$ , $-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)$ , $-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3

$-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)$ , $-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)$ , $-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3

$-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)$ , $-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)$ , $-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3

$-(b_{32}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,)}(G_{35},t)$, $-(b_{15}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$ are sixth detrition coefficients for category 1, 2 and 3

$-(b_{13}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{38}^{''})^{(7,7,)}(G_{39},t)$ are seventh detrition coefficients for category 1, 2 and 3

$-(b_{36}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{37}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)$, $-(b_{38}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)$ are eighth detrition coefficients for category 1, 2 and 3

$-(b_{44}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{45}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$, $-(b_{46}^{''})^{(9,9,9,9,9,9,9,9,9)}(G_{47},t)$ are ninth detrition coefficients for category 1, 2 and 3

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$$\frac{dG_{44}}{dt} = (a_{44})^{(9)}G_{45} - \begin{bmatrix} (a_{44}^{'})^{(9)}\left[+(a_{44}^{''})^{(9)}(T_{45},t)\right] & \left[+(a_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)\right] & \left[+(a_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)\right] \\ \left[+(a_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)\right] & \left[+(a_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)\right] & \left[+(a_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)\right] \\ \left[+(a_{13}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)\right] & \left[+(a_{36}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)\right] & \left[+(a_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)\right] \end{bmatrix} G_{13}$$

$$\frac{dG_{45}}{dt} = (a_{45})^{(9)}G_{44} - \begin{bmatrix} (a_{45}^{'})^{(9)}\left[+(a_{45}^{''})^{(9)}(T_{45},t)\right] & \left[+(a_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)\right] & \left[+(a_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)\right] \\ \left[+(a_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)\right] & \left[+(a_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)\right] & \left[+(a_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)\right] \\ \left[+(a_{14}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)\right] & \left[+(a_{37}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)\right] & \left[+(a_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)\right] \end{bmatrix} G_{14}$$

$$\frac{dG_{46}}{dt} = (a_{46})^{(9)}G_{45} - \begin{bmatrix} (a_{46}^{'})^{(9)}\left[+(a_{46}^{''})^{(9)}(T_{37},t)\right] & \left[+(a_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)\right] & \left[+(a_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)\right] \\ \left[+(a_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)\right] & \left[+(a_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)\right] & \left[+(a_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)\right] \\ \left[+(a_{15}^{''})^{(1,1,1,1,1,1,1,1,1)}(T_{14},t)\right] & \left[+(a_{38}^{''})^{(7,7,7,7,7,7,7,7,7)}(T_{37},t)\right] & \left[+(a_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(T_{41},t)\right] \end{bmatrix} G_{15}$$

Where $+(a_{44}^{''})^{(9)}(T_{45},t)$, $+(a_{45}^{''})^{(9)}(T_{45},t)$, $+(a_{46}^{''})^{(9)}(T_{37},t)$ are first augmentation coefficients for category 1, 2 and 3

$+(a_{16}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{17}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)$, $+(a_{18}^{''})^{(2,2,2,2,2,2,2,2,2)}(T_{17},t)$ are second augmentation coefficient for category 1, 2 and 3

$+(a_{20}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{21}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)$, $+(a_{22}^{''})^{(3,3,3,3,3,3,3,3,3)}(T_{21},t)$ are third augmentation coefficient for category 1, 2 and 3

$+(a_{24}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{25}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$, $+(a_{26}^{''})^{(4,4,4,4,4,4,4,4,4)}(T_{25},t)$ are fourth augmentation coefficient for category 1, 2 and 3

$+(a_{28}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{29}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$, $+(a_{30}^{''})^{(5,5,5,5,5,5,5,5,5)}(T_{29},t)$ are fifth augmentation coefficient for category 1, 2 and 3

$+(a_{32}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{33}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$, $+(a_{34}^{''})^{(6,6,6,6,6,6,6,6,6)}(T_{33},t)$ are sixth

augmentation coefficient for category 1, 2 and 3
$+(a_{13}^{''})^{(1,1,1,1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{14}^{''})^{(1,1,1,1,1,1,1,1,1,1)}(T_{14},t)$, $+(a_{15}^{''})^{(1,1,1,1,1,1,1,1,1,1)}(T_{14},t)$ are Seventh augmentation coefficient for category 1, 2 and 3
$+(a_{38}^{''})^{(7,7,7,7,7,7,7,7,7,7)}(T_{37},t)$ $+(a_{37}^{''})^{(7,7,7,7,7,7,7,7,7,7)}(T_{37},t)$ $+(a_{36}^{''})^{(7,7,7,7,7,7,7,7,7,7)}(T_{37},t)$ are eighth augmentation coefficient for 1,2,3
$+(a_{40}^{''})^{(8,8,8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{42}^{''})^{(8,8,8,8,8,8,8,8,8,8)}(T_{41},t)$, $+(a_{41}^{''})^{(8,8,8,8,8,8,8,8,8,8)}(T_{41},t)$ are ninth augmentation coefficient for 1,2,3

$$\frac{dT_{44}}{dt} =$$
$$(b_{44})^{(9)}T_{45} -$$
$$\begin{bmatrix} (b_{44}^{'})^{(9)}\left[-(b_{44}^{''})^{(9)}(G_{47},t)\right] & \left[-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2,2)}(G_{19},t)\right] & \left[-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3,3)}(G_{23},t)\right] \\ \left[-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4,4)}(G_{27},t)\right] & \left[-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5,5)}(G_{31},t)\right] & \left[-(b_{32}^{''})^{(6,6,6,6,6,6,6,6,6,6)}(G_{35},t)\right] \\ \left[-(b_{13}^{''})^{(1,1,1,1,1,1,1,1,1,1)}(G,t)\right] & \left[-(b_{36}^{''})^{(7,7,7,7,7,7,7,7,7,7)}(G_{39},t)\right] & \left[-(b_{40}^{''})^{(8,8,8,8,8,8,8,8,8,8)}(G_{43},t)\right] \end{bmatrix} T_{13}$$

$$\frac{dT_{45}}{dt} =$$
$$(b_{45})^{(9)}T_{44} - \begin{bmatrix} (b_{45}^{'})^{(9)}\left[-(b_{45}^{''})^{(9)}(G_{47},t)\right] & \left[-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2,2)}(G_{19},t)\right] & \left[-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3,3)}(G_{23},t)\right] \\ \left[-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4,4)}(G_{27},t)\right] & \left[-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5,5)}(G_{31},t)\right] & \left[-(b_{33}^{''})^{(6,6,6,6,6,6,6,6,6,6)}(G_{35},t)\right] \\ \left[-(b_{14}^{''})^{(1,1,1,1,1,1,1,1,1,1)}(G,t)\right] & \left[-(b_{37}^{''})^{(7,7,7,7,7,7,7,7,7,7)}(G_{39},t)\right] & \left[-(b_{41}^{''})^{(8,8,8,8,8,8,8,8,8,8)}(G_{43},t)\right] \end{bmatrix} T_{14}$$

$$\frac{dT_{46}}{dt} =$$
$$(b_{46})^{(9)}T_{45} - \begin{bmatrix} (b_{46}^{'})^{(9)}\left[-(b_{46}^{''})^{(9)}(G_{47},t)\right] & \left[-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2,2)}(G_{19},t)\right] & \left[-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3,3)}(G_{23},t)\right] \\ \left[-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4,4)}(G_{27},t)\right] & \left[-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5,5)}(G_{31},t)\right] & \left[-(b_{34}^{''})^{(6,6,6,6,6,6,6,6,6,6)}(G_{35},t)\right] \\ \left[-(b_{15}^{''})^{(1,1,1,1,1,1,1,1,1,1)}(G,t)\right] & \left[-(b_{38}^{''})^{(7,7,7,7,7,7,7,7,7,7)}(G_{39},t)\right] & \left[-(b_{42}^{''})^{(8,8,8,8,8,8,8,8,8,8)}(G_{43},t)\right] \end{bmatrix} T_{15}$$

Where $-(b_{44}^{''})^{(9)}(G_{47},t)$, $-(b_{45}^{''})^{(9)}(G_{47},t)$, $-(b_{46}^{''})^{(9)}(G_{47},t)$ are first detrition coefficients for category 1, 2 and 3
$-(b_{16}^{''})^{(2,2,2,2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{17}^{''})^{(2,2,2,2,2,2,2,2,2,2)}(G_{19},t)$, $-(b_{18}^{''})^{(2,2,2,2,2,2,2,2,2,2)}(G_{19},t)$ are second detrition coefficients for category 1, 2 and 3
$-(b_{20}^{''})^{(3,3,3,3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{21}^{''})^{(3,3,3,3,3,3,3,3,3,3)}(G_{23},t)$, $-(b_{22}^{''})^{(3,3,3,3,3,3,3,3,3,3)}(G_{23},t)$ are third detrition coefficients for category 1, 2 and 3
$-(b_{24}^{''})^{(4,4,4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{25}^{''})^{(4,4,4,4,4,4,4,4,4,4)}(G_{27},t)$, $-(b_{26}^{''})^{(4,4,4,4,4,4,4,4,4,4)}(G_{27},t)$ are fourth detrition coefficients for category 1, 2 and 3
$-(b_{28}^{''})^{(5,5,5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{29}^{''})^{(5,5,5,5,5,5,5,5,5,5)}(G_{31},t)$, $-(b_{30}^{''})^{(5,5,5,5,5,5,5,5,5,5)}(G_{31},t)$ are fifth detrition coefficients for category 1, 2 and 3
$-(b_{32}^{''})^{(6,6,6,6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{33}^{''})^{(6,6,6,6,6,6,6,6,6,6)}(G_{35},t)$, $-(b_{34}^{''})^{(6,6,6,6,6,6,6,6,6,6)}(G_{35},t)$ are sixth detrition coefficients for category 1, 2 and 3
$-(b_{15}^{''})^{(1,1,1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{14}^{''})^{(1,1,1,1,1,1,1,1,1,1)}(G,t)$, $-(b_{13}^{''})^{(1,1,1,1,1,1,1,1,1,1)}(G,t)$ are seventh detrition coefficients for category 1, 2 and 3

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| $\boxed{-(b_{37}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{36}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$, $\boxed{-(b_{38}^{''})^{(7,7,7,7,7,7,7,7)}(G_{39},t)}$ are eighth detrition coefficients for category 1, 2 and 3<br>$\boxed{-(b_{42}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{41}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$, $\boxed{-(b_{40}^{''})^{(8,8,8,8,8,8,8,8,8)}(G_{43},t)}$ are ninth detrition coefficients for category 1, 2 and 3 | |
| Where we suppose | |
| $(a_i)^{(1)}, (a_i^{'})^{(1)}, (a_i^{''})^{(1)}, (b_i)^{(1)}, (b_i^{'})^{(1)}, (b_i^{''})^{(1)} > 0, \quad i,j = 13,14,15$<br><br>The functions $(a_i^{''})^{(1)}, (b_i^{''})^{(1)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(1)}$, $(r_i)^{(1)}$:<br><br>$(a_i^{''})^{(1)}(T_{14},t) \leq (p_i)^{(1)} \leq (\hat{A}_{13})^{(1)}$<br><br>$(b_i^{''})^{(1)}(G,t) \leq (r_i)^{(1)} \leq (b_i^{'})^{(1)} \leq (\hat{B}_{13})^{(1)}$ | *97* |
| $lim_{T_2\to\infty}(a_i^{''})^{(1)}(T_{14},t) = (p_i)^{(1)}$<br><br>$\lim_{G\to\infty}(b_i^{''})^{(1)}(G,t) = (r_i)^{(1)}$<br><br>**Definition of** $(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}$ :<br><br>Where $\boxed{(\hat{A}_{13})^{(1)}, (\hat{B}_{13})^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}}$ are positive constants and $\boxed{i = 13,14,15}$ | 98 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(1)}(T_{14}^{'},t) - (a_i^{''})^{(1)}(T_{14},t)\rvert \leq (\hat{k}_{13})^{(1)}\lvert T_{14} - T_{14}^{'}\rvert e^{-(\hat{M}_{13})^{(1)}t}$<br><br>$\lvert(b_i^{''})^{(1)}(G^{'},t) - (b_i^{''})^{(1)}(G,t)\rvert < (\hat{k}_{13})^{(1)}\lVert G - G^{'}\rVert e^{-(\hat{M}_{13})^{(1)}t}$ | 99 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(1)}(T_{14}^{'},t)$ and $(a_i^{''})^{(1)}(T_{14},t)$ . $(T_{14}^{'},t)$ and $(T_{14},t)$ are points belonging to the interval $\left[(\hat{k}_{13})^{(1)}, (\hat{M}_{13})^{(1)}\right]$. It is to be noted that $(a_i^{''})^{(1)}(T_{14},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{13})^{(1)} = 1$ then the function $(a_i^{''})^{(1)}(T_{14},t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}$ :<br><br>$(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}$, are positive constants<br><br>$\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}} < 1$ | 100 |
| **Definition of** $(\hat{P}_{13})^{(1)}, (\hat{Q}_{13})^{(1)}$ :<br><br>There exists two constants $(\hat{P}_{13})^{(1)}$ and $(\hat{Q}_{13})^{(1)}$ which together With $(\hat{M}_{13})^{(1)}, (\hat{k}_{13})^{(1)}, (\hat{A}_{13})^{(1)}$ and $(\hat{B}_{13})^{(1)}$ and the constants $(a_i)^{(1)}, (a_i^{'})^{(1)}, (b_i)^{(1)}, (b_i^{'})^{(1)}, (p_i)^{(1)}, (r_i)^{(1)}, i = 13,14,15,$ | 101 |

| | |
|---|---|
| satisfy the inequalities<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(a_i)^{(1)} + (a_i')^{(1)} + (\hat{A}_{13})^{(1)} + (\hat{P}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$<br>$\frac{1}{(\hat{M}_{13})^{(1)}}[(b_i)^{(1)} + (b_i')^{(1)} + (\hat{B}_{13})^{(1)} + (\hat{Q}_{13})^{(1)}(\hat{k}_{13})^{(1)}] < 1$ | |
| Where we suppose | |
| $(a_i)^{(2)}, (a_i')^{(2)}, (a_i'')^{(2)}, (b_i)^{(2)}, (b_i')^{(2)}, (b_i'')^{(2)} > 0, \quad i,j = 16,17,18$ | |
| The functions $(a_i'')^{(2)}, (b_i'')^{(2)}$ are positive continuous increasing and bounded. | |
| **Definition of** $(p_i)^{(2)}$, $(r_i)^{(2)}$: | |
| $(a_i'')^{(2)}(T_{17}, t) \leq (p_i)^{(2)} \leq (\hat{A}_{16})^{(2)}$ | 102 |
| $(b_i'')^{(2)}(G_{19}, t) \leq (r_i)^{(2)} \leq (b_i')^{(2)} \leq (\hat{B}_{16})^{(2)}$ | 103 |
| $\lim_{T_2 \to \infty} (a_i'')^{(2)}(T_{17}, t) = (p_i)^{(2)}$ | 104 |
| $\lim_{G \to \infty} (b_i'')^{(2)}((G_{19}), t) = (r_i)^{(2)}$ | 105 |
| **Definition of** $(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}$ :<br>Where $\boxed{(\hat{A}_{16})^{(2)}, (\hat{B}_{16})^{(2)}, (p_i)^{(2)}, (r_i)^{(2)}}$ are positive constants and $\boxed{i = 16,17,18}$ | 106 |
| They satisfy Lipschitz condition: | |
| $\lvert(a_i'')^{(2)}(T_{17}', t) - (a_i'')^{(2)}(T_{17}, t)\rvert \leq (\hat{k}_{16})^{(2)} \lvert T_{17} - T_{17}'\rvert e^{-(\hat{M}_{16})^{(2)}t}$ | 107 |
| $\lvert(b_i'')^{(2)}((G_{19})', t) - (b_i'')^{(2)}((G_{19}), t)\rvert < (\hat{k}_{16})^{(2)} \lvert\lvert(G_{19}) - (G_{19})'\rvert\rvert e^{-(\hat{M}_{16})^{(2)}t}$ | 108 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(2)}(T_{17}', t)$ and $(a_i'')^{(2)}(T_{17}, t)$ . $(T_{17}', t)$ and $(T_{17}, t)$ are points belonging to the interval $[(\hat{k}_{16})^{(2)}, (\hat{M}_{16})^{(2)}]$ . It is to be noted that $(a_i'')^{(2)}(T_{17}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{16})^{(2)} = 1$ then the function $(a_i'')^{(2)}(T_{17}, t)$ , the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$ : | |
| $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}$, are positive constants<br>$\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}, \frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ | 109 |
| **Definition of** $(\hat{P}_{13})^{(2)}, (\hat{Q}_{13})^{(2)}$ :<br>There exists two constants $(\hat{P}_{16})^{(2)}$ and $(\hat{Q}_{16})^{(2)}$ which together | |

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| with $(\hat{M}_{16})^{(2)}, (\hat{k}_{16})^{(2)}, (\hat{A}_{16})^{(2)} and\ (\hat{B}_{16})^{(2)}$ and the constants $(a_i)^{(2)}, (a_i^{'})^{(2)}, (b_i)^{(2)}, (b_i^{'})^{(2)}, (p_i)^{(2)}, (r_i)^{(2)}, i = 16,17,18,$ satisfy the inequalities | |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(a_i)^{(2)} + (a_i^{'})^{(2)} + (\hat{A}_{16})^{(2)} + (\hat{P}_{16})^{(2)} (\hat{k}_{16})^{(2)}] < 1$ | 110 |
| $\frac{1}{(\hat{M}_{16})^{(2)}}[(b_i)^{(2)} + (b_i^{'})^{(2)} + (\hat{B}_{16})^{(2)} + (\hat{Q}_{16})^{(2)} (\hat{k}_{16})^{(2)}] < 1$ | 111 |
| Where we suppose | |
| $(a_i)^{(3)}, (a_i^{'})^{(3)}, (a_i^{''})^{(3)}, (b_i)^{(3)}, (b_i^{'})^{(3)}, (b_i^{''})^{(3)} > 0, \quad i, j = 20,21,22$<br>The functions $(a_i^{''})^{(3)}, (b_i^{''})^{(3)}$ are positive continuous increasing and bounded.<br>**Definition of** $(p_i)^{(3)}$, $(r_i)^{(3)}$:<br>$(a_i^{''})^{(3)}(T_{21}, t) \leq (p_i)^{(3)} \leq (\hat{A}_{20})^{(3)}$<br>$(b_i^{''})^{(3)}(G_{23}, t) \leq (r_i)^{(3)} \leq (b_i^{'})^{(3)} \leq (\hat{B}_{20})^{(3)}$ | 112 |
| $lim_{T_2 \to \infty}(a_i^{''})^{(3)}(T_{21}, t) = (p_i)^{(3)}$<br>$\lim_{G \to \infty}(b_i^{''})^{(3)}(G_{23}, t) = (r_i)^{(3)}$<br>**Definition of** $(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}$ :<br>Where $\boxed{(\hat{A}_{20})^{(3)}, (\hat{B}_{20})^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}}$ are positive constants and $\boxed{i = 20,21,22}$ | 113 |
| They satisfy Lipschitz condition:<br>$\lvert(a_i^{''})^{(3)}(T_{21}^{'}, t) - (a_i^{''})^{(3)}(T_{21}, t)\rvert \leq (\hat{k}_{20})^{(3)} \lvert T_{21} - T_{21}^{'}\rvert e^{-(\hat{M}_{20})^{(3)}t}$<br>$\lvert(b_i^{''})^{(3)}(G_{23}^{'}, t) - (b_i^{''})^{(3)}(G_{23}, t)\rvert < (\hat{k}_{20})^{(3)} \lvert\lvert G_{23} - G_{23}^{'}\rvert\rvert e^{-(\hat{M}_{20})^{(3)}t}$ | 114 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(3)}(T_{21}^{'}, t)$ and $(a_i^{''})^{(3)}(T_{21}, t)$ . $(T_{21}^{'}, t)$ And $(T_{21}, t)$ are points belonging to the interval $[(\hat{k}_{20})^{(3)}, (\hat{M}_{20})^{(3)}]$. It is to be noted that $(a_i^{''})^{(3)}(T_{21}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{20})^{(3)} = 1$ then the function $(a_i^{''})^{(3)}(T_{21}, t)$, the first augmentation coefficient attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$ :<br>$(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}$, are positive constants<br>$\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}, \frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}} < 1$ | 115 |

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| There exists two constants There exists two constants $(\hat{P}_{20})^{(3)}$ and $(\hat{Q}_{20})^{(3)}$ which together with $(\hat{M}_{20})^{(3)}, (\hat{k}_{20})^{(3)}, (\hat{A}_{20})^{(3)} and\ (\hat{B}_{20})^{(3)}$ and the constants $(a_i)^{(3)}, (a_i^{'})^{(3)}, (b_i)^{(3)}, (b_i^{'})^{(3)}, (p_i)^{(3)}, (r_i)^{(3)}, i = 20,21,22,$ satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(a_i)^{(3)} + (a_i^{'})^{(3)} + (\hat{A}_{20})^{(3)} + (\hat{P}_{20})^{(3)} (\hat{k}_{20})^{(3)}] < 1$<br><br>$\frac{1}{(\hat{M}_{20})^{(3)}}[(b_i)^{(3)} + (b_i^{'})^{(3)} + (\hat{B}_{20})^{(3)} + (\hat{Q}_{20})^{(3)} (\hat{k}_{20})^{(3)}] < 1$ | 116 |
| Where we suppose | |
| $(a_i)^{(4)}, (a_i^{'})^{(4)}, (a_i^{''})^{(4)}, (b_i)^{(4)}, (b_i^{'})^{(4)}, (b_i^{''})^{(4)} > 0, \quad i,j = 24,25,26$<br><br>The functions $(a_i^{''})^{(4)}, (b_i^{''})^{(4)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(4)}$, $(r_i)^{(4)}$:<br><br>$(a_i^{''})^{(4)}(T_{25}, t) \leq (p_i)^{(4)} \leq (\hat{A}_{24})^{(4)}$<br><br>$(b_i^{''})^{(4)}((G_{27}), t) \leq (r_i)^{(4)} \leq (b_i^{'})^{(4)} \leq (\hat{B}_{24})^{(4)}$ | 117 |
| $lim_{T_2 \to \infty} (a_i^{''})^{(4)}(T_{25}, t) = (p_i)^{(4)}$<br><br>$\lim_{G \to \infty} (b_i^{''})^{(4)}((G_{27}), t) = (r_i)^{(4)}$<br><br>**Definition of** $(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}$ :<br><br>Where $\boxed{(\hat{A}_{24})^{(4)}, (\hat{B}_{24})^{(4)}, (p_i)^{(4)}, (r_i)^{(4)}}$ are positive constants and $\boxed{i = 24,25,26}$ | 118 |
| They satisfy Lipschitz condition:<br><br>$\|(a_i^{''})^{(4)}(T_{25}^{'}, t) - (a_i^{''})^{(4)}(T_{25}, t)\| \leq (\hat{k}_{24})^{(4)} \|T_{25} - T_{25}^{'}\| e^{-(\hat{M}_{24})^{(4)} t}$<br><br>$\|(b_i^{''})^{(4)}((G_{27})', t) - (b_i^{''})^{(4)}((G_{27}), t)\| < (\hat{k}_{24})^{(4)} \|(G_{27}) - (G_{27})'\| e^{-(\hat{M}_{24})^{(4)} t}$ | 119 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(4)}(T_{25}^{'}, t)$ and $(a_i^{''})^{(4)}(T_{25}, t)$ . $(T_{25}^{'}, t)$ and $(T_{25}, t)$ are points belonging to the interval $[(\hat{k}_{24})^{(4)}, (\hat{M}_{24})^{(4)}]$. It is to be noted that $(a_i^{''})^{(4)}(T_{25}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{24})^{(4)} = 1$ then the function $(a_i^{''})^{(4)}(T_{25}, t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$ :<br><br>$(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}$, are positive constants | 120 |

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| $\frac{(a_i)^{(4)}}{(\hat{M}_{24})^{(4)}}, \frac{(b_i)^{(4)}}{(\hat{M}_{24})^{(4)}} < 1$ | |
| **Definition of** $(\hat{P}_{24})^{(4)}, (\hat{Q}_{24})^{(4)}$ :<br><br>There exists two constants $(\hat{P}_{24})^{(4)}$ and $(\hat{Q}_{24})^{(4)}$ which together with $(\hat{M}_{24})^{(4)}, (\hat{k}_{24})^{(4)}, (\hat{A}_{24})^{(4)} and\ (\hat{B}_{24})^{(4)}$ and the constants $(a_i)^{(4)}, (a_i^{'})^{(4)}, (b_i)^{(4)}, (b_i^{'})^{(4)}, (p_i)^{(4)}, (r_i)^{(4)}, i = 24,25,26$, satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{24})^{(4)}}[(a_i)^{(4)} + (a_i^{'})^{(4)} + (\hat{A}_{24})^{(4)} + (\hat{P}_{24})^{(4)} (\hat{k}_{24})^{(4)}] < 1$<br><br>$\frac{1}{(\hat{M}_{24})^{(4)}}[(b_i)^{(4)} + (b_i^{'})^{(4)} + (\hat{B}_{24})^{(4)} + (\hat{Q}_{24})^{(4)} (\hat{k}_{24})^{(4)}] < 1$ | 121 |
| Where we suppose | |
| $(a_i)^{(5)}, (a_i^{'})^{(5)}, (a_i^{''})^{(5)}, (b_i)^{(5)}, (b_i^{'})^{(5)}, (b_i^{''})^{(5)} > 0, \quad i,j = 28,29,30$<br><br>The functions $(a_i^{''})^{(5)}, (b_i^{''})^{(5)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(5)}, (r_i)^{(5)}$:<br><br>$(a_i^{''})^{(5)}(T_{29}, t) \leq (p_i)^{(5)} \leq (\hat{A}_{28})^{(5)}$<br><br>$(b_i^{''})^{(5)}((G_{31}), t) \leq (r_i)^{(5)} \leq (b_i^{'})^{(5)} \leq (\hat{B}_{28})^{(5)}$ | 122 |
| $lim_{T_2 \to \infty} (a_i^{''})^{(5)} (T_{29}, t) = (p_i)^{(5)}$<br><br>$\lim_{G \to \infty} (b_i^{''})^{(5)} (G_{31}, t) = (r_i)^{(5)}$<br><br>**Definition of** $(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}$ :<br><br>Where $\boxed{(\hat{A}_{28})^{(5)}, (\hat{B}_{28})^{(5)}, (p_i)^{(5)}, (r_i)^{(5)}}$ are positive constants and $\boxed{i = 28,29,30}$ | 123 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i^{''})^{(5)}(T_{29}^{'}, t) - (a_i^{''})^{(5)}(T_{29}, t)\rvert \leq (\hat{k}_{28})^{(5)} \lvert T_{29} - T_{29}^{'}\rvert e^{-(\hat{M}_{28})^{(5)} t}$<br><br>$\lvert(b_i^{''})^{(5)}((G_{31})', t) - (b_i^{''})^{(5)}((G_{31}), t)\rvert < (\hat{k}_{28})^{(5)} \lvert\lvert(G_{31}) - (G_{31})'\rvert\rvert e^{-(\hat{M}_{28})^{(5)} t}$ | 124 |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(5)}(T_{29}^{'}, t)$ and $(a_i^{''})^{(5)}(T_{29}, t)$ . $(T_{29}^{'}, t)$ and $(T_{29}, t)$ are points belonging to the interval $[(\hat{k}_{28})^{(5)}, (\hat{M}_{28})^{(5)}]$. It is to be noted that $(a_i^{''})^{(5)}(T_{29}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{28})^{(5)} = 1$ then the function $(a_i^{''})^{(5)}(T_{29}, t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}$ : | 125 |

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| $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}$, are positive constants<br><br>$\frac{(a_i)^{(5)}}{(\hat{M}_{28})^{(5)}}, \frac{(b_i)^{(5)}}{(\hat{M}_{28})^{(5)}} < 1$ | |
| **Definition of** $(\hat{P}_{28})^{(5)}, (\hat{Q}_{28})^{(5)}$ :<br><br>There exists two constants $(\hat{P}_{28})^{(5)}$ and $(\hat{Q}_{28})^{(5)}$ which together with $(\hat{M}_{28})^{(5)}, (\hat{k}_{28})^{(5)}, (\hat{A}_{28})^{(5)} and\ (\hat{B}_{28})^{(5)}$ and the constants $(a_i)^{(5)}, (a_i')^{(5)}, (b_i)^{(5)}, (b_i')^{(5)}, (p_i)^{(5)}, (r_i)^{(5)}, i = 28,29,30,$ satisfy the inequalities<br><br>$\frac{1}{(\hat{M}_{28})^{(5)}}[(a_i)^{(5)} + (a_i')^{(5)} + (\hat{A}_{28})^{(5)} + (\hat{P}_{28})^{(5)} (\hat{k}_{28})^{(5)}] < 1$<br><br>$\frac{1}{(\hat{M}_{28})^{(5)}}[(b_i)^{(5)} + (b_i')^{(5)} + (\hat{B}_{28})^{(5)} + (\hat{Q}_{28})^{(5)} (\hat{k}_{28})^{(5)}] < 1$ | 126 |
| Where we suppose | |
| $(a_i)^{(6)}, (a_i')^{(6)}, (a_i'')^{(6)}, (b_i)^{(6)}, (b_i')^{(6)}, (b_i'')^{(6)} > 0, \quad i, j = 32,33,34$<br><br>The functions $(a_i'')^{(6)}, (b_i'')^{(6)}$ are positive continuous increasing and bounded.<br><br>**Definition of** $(p_i)^{(6)}$, $(r_i)^{(6)}$:<br><br>$(a_i'')^{(6)}(T_{33}, t) \le (p_i)^{(6)} \le (\hat{A}_{32})^{(6)}$<br><br>$(b_i'')^{(6)}((G_{35}), t) \le (r_i)^{(6)} \le (b_i')^{(6)} \le (\hat{B}_{32})^{(6)}$ | 127 |
| $lim_{T_2 \to \infty} (a_i'')^{(6)} (T_{33}, t) = (p_i)^{(6)}$<br><br>$\lim_{G \to \infty} (b_i'')^{(6)} ((G_{35}), t) = (r_i)^{(6)}$<br><br>**Definition of** $(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}$ :<br><br>Where $\boxed{(\hat{A}_{32})^{(6)}, (\hat{B}_{32})^{(6)}, (p_i)^{(6)}, (r_i)^{(6)}}$ are positive constants and $\boxed{i = 32,33,34}$ | 128 |
| They satisfy Lipschitz condition:<br><br>$\lvert(a_i'')^{(6)}(T_{33}', t) - (a_i'')^{(6)}(T_{33}, t)\rvert \le (\hat{k}_{32})^{(6)} \lvert T_{33} - T_{33}'\rvert e^{-(\hat{M}_{32})^{(6)} t}$<br><br>$\lvert(b_i'')^{(6)}((G_{35})', t) - (b_i'')^{(6)}((G_{35}), t)\rvert < (\hat{k}_{32})^{(6)} \lvert\lvert(G_{35}) - (G_{35})'\rvert\rvert e^{-(\hat{M}_{32})^{(6)} t}$ | |
| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(6)}(T_{33}', t)$ and $(a_i'')^{(6)}(T_{33}, t)$ . $(T_{33}', t)$ and $(T_{33}, t)$ are points belonging to the interval $[(\hat{k}_{32})^{(6)}, (\hat{M}_{32})^{(6)}]$. It is to be noted that $(a_i'')^{(6)}(T_{33}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{32})^{(6)} = 1$ then the function $(a_i'')^{(6)}(T_{33}, t)$, the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |

**Definition of** $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}$ : 129

$(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)},$ are positive constants

$$\frac{(a_i)^{(6)}}{(\hat{M}_{32})^{(6)}}, \frac{(b_i)^{(6)}}{(\hat{M}_{32})^{(6)}} < 1$$

**Definition of** $(\hat{P}_{32})^{(6)}, (\hat{Q}_{32})^{(6)}$ : 130

There exists two constants $(\hat{P}_{32})^{(6)}$ and $(\hat{Q}_{32})^{(6)}$ which together with $(\hat{M}_{32})^{(6)}, (\hat{k}_{32})^{(6)}, (\hat{A}_{32})^{(6)} and\ (\hat{B}_{32})^{(6)}$ and the constants $(a_i)^{(6)}, (a_i')^{(6)}, (b_i)^{(6)}, (b_i')^{(6)}, (p_i)^{(6)}, (r_i)^{(6)}, i = 32,33,34,$ satisfy the inequalities

$$\frac{1}{(\hat{M}_{32})^{(6)}}[(a_i)^{(6)} + (a_i')^{(6)} + (\hat{A}_{32})^{(6)} + (\hat{P}_{32})^{(6)}(\hat{k}_{32})^{(6)}] < 1$$

$$\frac{1}{(\hat{M}_{32})^{(6)}}[(b_i)^{(6)} + (b_i')^{(6)} + (\hat{B}_{32})^{(6)} + (\hat{Q}_{32})^{(6)}(\hat{k}_{32})^{(6)}] < 1$$

Where we suppose

(QQ) $(a_i)^{(7)}, (a_i')^{(7)}, (a_i'')^{(7)}, (b_i)^{(7)}, (b_i')^{(7)}, (b_i'')^{(7)} > 0, \quad i, j = 36,37,38$ 131

(RR) The functions $(a_i'')^{(7)}, (b_i'')^{(7)}$ are positive continuous increasing and bounded.

**Definition of** $(p_i)^{(7)}$, $(r_i)^{(7)}$:

$$(a_i'')^{(7)}(T_{37}, t) \le (p_i)^{(7)} \le (\hat{A}_{36})^{(7)}$$

$$(b_i'')^{(7)}(G_{39}, t) \le (r_i)^{(7)} \le (b_i')^{(7)} \le (\hat{B}_{36})^{(7)}$$

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*(SS)* $\lim_{T_2 \to \infty} (a_i'')^{(7)}(T_{37}, t) = (p_i)^{(7)}$

*(TT)*

$\lim_{G \to \infty} (b_i'')^{(7)}((G_{39}), t) = (r_i)^{(7)}$

**Definition of** $(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}$ :

Where $\boxed{(\hat{A}_{36})^{(7)}, (\hat{B}_{36})^{(7)}, (p_i)^{(7)}, (r_i)^{(7)}}$ are positive constants and $\boxed{i = 36,37,38}$

They satisfy Lipschitz condition: 133

$$|(a_i'')^{(7)}(T_{37}', t) - (a_i'')^{(7)}(T_{37}, t)| \le (\hat{k}_{36})^{(7)} |T_{37} - T_{37}'| e^{-(\hat{M}_{36})^{(7)} t}$$

$$|(b_i'')^{(7)}((G_{39})', t) - (b_i'')^{(7)}((G_{39}), t)| < (\hat{k}_{36})^{(7)} ||(G_{39}) - (G_{39})'|| e^{-(\hat{M}_{36})^{(7)} t}$$

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| With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(7)}(T_{37}^{'},t)$ and$(a_i^{''})^{(7)}(T_{37},t)$ . $(T_{37}^{'},t)$ and $(T_{37},t)$ are points belonging to the interval $\left[(\hat{k}_{36})^{(7)},(\widehat{M}_{36})^{(7)}\right]$. It is to be noted that $(a_i^{''})^{(7)}(T_{37},t)$ is uniformly continuous. In the eventuality of the fact, that if $(\widehat{M}_{36})^{(7)}=1$ then the function $(a_i^{''})^{(7)}(T_{37},t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous. | |
| **Definition of** $(\widehat{M}_{36})^{(7)},(\hat{k}_{36})^{(7)}$ : <br><br>(UU) $(\widehat{M}_{36})^{(7)},(\hat{k}_{36})^{(7)}$, are positive constants <br><br>$\frac{(a_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}$ , $\frac{(b_i)^{(7)}}{(\widehat{M}_{36})^{(7)}}<1$ | 134 |
| **Definition of** $(\hat{P}_{36})^{(7)},(\hat{Q}_{36})^{(7)}$ : <br><br>(VV) There exists two constants $(\hat{P}_{36})^{(7)}$ and $(\hat{Q}_{36})^{(7)}$ which together with $(\widehat{M}_{36})^{(7)},(\hat{k}_{36})^{(7)},(\hat{A}_{36})^{(7)} and\ (\hat{B}_{36})^{(7)}$ and the constants $(a_i)^{(7)},(a_i^{'})^{(7)},(b_i)^{(7)},(b_i^{'})^{(7)},(p_i)^{(7)},\ (r_i)^{(7)}, i=36,37,38$, satisfy the inequalities <br><br>$\frac{1}{(\widehat{M}_{36})^{(7)}}[(a_i)^{(7)}+(a_i^{'})^{(7)}+(\hat{A}_{36})^{(7)}+(\hat{P}_{36})^{(7)}(\hat{k}_{36})^{(7)}]<1$ <br><br>$\frac{1}{(\widehat{M}_{36})^{(7)}}[(b_i)^{(7)}+(b_i^{'})^{(7)}+(\hat{B}_{36})^{(7)}+(\hat{Q}_{36})^{(7)}(\hat{k}_{36})^{(7)}]<1$ | 135 |
| Where we suppose | |
| $(a_i)^{(8)},(a_i^{'})^{(8)},(a_i^{''})^{(8)},(b_i)^{(8)},(b_i^{'})^{(8)},(b_i^{''})^{(8)}>0, \quad i,j=40,41,42$ | 136 |
| The functions $(a_i^{''})^{(8)},(b_i^{''})^{(8)}$ are positive continuous increasing and bounded | |
| **Definition of** $(p_i)^{(8)},\ (r_i)^{(8)}$: | 137 |
| $(a_i^{''})^{(8)}(T_{41},t)\leq(p_i)^{(8)}\leq(\hat{A}_{40})^{(8)}$ | 138 |
| $(b_i^{''})^{(8)}((G_{43}),t)\leq\ (r_i)^{(8)}\leq(b_i^{'})^{(8)}\leq(\hat{B}_{40})^{(8)}$ | 139 |
| $\lim_{T_2\to\infty}(a_i^{''})^{(8)}(T_{41},t)=(p_i)^{(8)}$ | 140 |
| $\lim_{G\to\infty}(b_i^{''})^{(8)}\left((G_{43}),t\right)=\ (r_i)^{(8)}$ | 141 |
| **Definition of** $(\hat{A}_{40})^{(8)},(\hat{B}_{40})^{(8)}$ : <br><br>Where $\boxed{(\hat{A}_{40})^{(8)},(\hat{B}_{40})^{(8)},(p_i)^{(8)},\ (r_i)^{(8)}}$ are positive constants and $\boxed{i=40,41,42}$ | |

They satisfy Lipschitz condition:

$$|(a_i'')^{(8)}(T_{41}', t) - (a_i'')^{(8)}(T_{41}, t)| \leq (\hat{k}_{40})^{(8)}|T_{41} - T_{41}'|e^{-(\hat{M}_{40})^{(8)}t} \quad (142)$$

$$|(b_i'')^{(8)}((G_{43})', t) - (b_i'')^{(8)}((G_{43}), t)| < (\hat{k}_{40})^{(8)}||(G_{43}) - (G_{43})'||e^{-(\hat{M}_{40})^{(8)}t} \quad (143)$$

With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i'')^{(8)}(T_{41}', t)$ and $(a_i'')^{(8)}(T_{41}, t)$ . $(T_{41}', t)$ and $(T_{41}, t)$ are points belonging to the interval $[(\hat{k}_{40})^{(8)}, (\hat{M}_{40})^{(8)}]$. It is to be noted that $(a_i'')^{(8)}(T_{41}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{40})^{(8)} = 1$ then the function $(a_i'')^{(8)}(T_{41}, t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

**Definition of** $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}$ :

$(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}$, are positive constants

$$\frac{(a_i)^{(8)}}{(\hat{M}_{40})^{(8)}}, \frac{(b_i)^{(8)}}{(\hat{M}_{40})^{(8)}} < 1 \quad (144)$$

**Definition of** $(\hat{P}_{40})^{(8)}, (\hat{Q}_{40})^{(8)}$ :

There exists two constants $(\hat{P}_{40})^{(8)}$ and $(\hat{Q}_{40})^{(8)}$ which together with $(\hat{M}_{40})^{(8)}, (\hat{k}_{40})^{(8)}, (\hat{A}_{40})^{(8)}$ $(\hat{B}_{40})^{(8)}$ and the constants $(a_i)^{(8)}, (a_i')^{(8)}, (b_i)^{(8)}, (b_i')^{(8)}, (p_i)^{(8)}, (r_i)^{(8)}, i = 40,41,42,$
Satisfy the inequalities

$$\frac{1}{(\hat{M}_{40})^{(8)}}[(a_i)^{(8)} + (a_i')^{(8)} + (\hat{A}_{40})^{(8)} + (\hat{P}_{40})^{(8)}(\hat{k}_{40})^{(8)}] < 1 \quad (145)$$

$$\frac{1}{(\hat{M}_{40})^{(8)}}[(b_i)^{(8)} + (b_i')^{(8)} + (\hat{B}_{40})^{(8)} + (\hat{Q}_{40})^{(8)}(\hat{k}_{40})^{(8)}] < 1 \quad (146)$$

Where we suppose

$$(a_i)^{(9)}, (a_i')^{(9)}, (a_i'')^{(9)}, (b_i)^{(9)}, (b_i')^{(9)}, (b_i'')^{(9)} > 0, \quad i, j = 44,45,46 \quad (146A)$$

The functions $(a_i'')^{(9)}, (b_i'')^{(9)}$ are positive continuous increasing and bounded.

**Definition of** $(p_i)^{(9)}$, $(r_i)^{(9)}$:

$$(a_i'')^{(9)}(T_{45}, t) \leq (p_i)^{(9)} \leq (\hat{A}_{44})^{(9)}$$

$$(b_i'')^{(9)}(G_{47}, t) \leq (r_i)^{(9)} \leq (b_i')^{(9)} \leq (\hat{B}_{44})^{(9)}$$

$$\lim_{T_2 \to \infty} (a_i'')^{(9)}(T_{45}, t) = (p_i)^{(9)}$$

$$\lim_{G \to \infty} (b_i'')^{(9)}(G_{47}, t) = (r_i)^{(9)}$$

**Definition of** $(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}$ :

Where $\boxed{(\hat{A}_{44})^{(9)}, (\hat{B}_{44})^{(9)}, (p_i)^{(9)}, (r_i)^{(9)}}$ are positive constants and $\boxed{i = 44,45,46}$

They satisfy Lipschitz condition:

$$|(a_i^{''})^{(9)}(T_{45}^{'}, t) - (a_i^{''})^{(9)}(T_{45}, t)| \leq (\hat{k}_{44})^{(9)}|T_{45} - T_{45}^{'}|e^{-(\hat{M}_{44})^{(9)}t}$$

$$|(b_i^{''})^{(9)}((G_{47})^{'}, t) - (b_i^{''})^{(9)}\big((G_{47}), t\big)| < (\hat{k}_{44})^{(9)}||(G_{47}) - (G_{47})^{'}||e^{-(\hat{M}_{44})^{(9)}t}$$

With the Lipschitz condition, we place a restriction on the behavior of functions $(a_i^{''})^{(9)}(T_{45}^{'}, t)$ and$(a_i^{''})^{(9)}(T_{45}, t)$ . $(T_{45}^{'}, t)$ and $(T_{45}, t)$ are points belonging to the interval $\left[(\hat{k}_{44})^{(9)}, (\hat{M}_{44})^{(9)}\right]$. It is to be noted that $(a_i^{''})^{(9)}(T_{45}, t)$ is uniformly continuous. In the eventuality of the fact, that if $(\hat{M}_{44})^{(9)} = 1$ then the function $(a_i^{''})^{(9)}(T_{45}, t)$ , the **first augmentation coefficient** attributable to the system, would be absolutely continuous.

**Definition of** $(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}$ :

$(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}$, are positive constants

$$\frac{(a_i)^{(9)}}{(\hat{M}_{44})^{(9)}} \ , \frac{(b_i)^{(9)}}{(\hat{M}_{44})^{(9)}} < 1$$

**Definition of** $(\hat{P}_{44})^{(9)}, (\hat{Q}_{44})^{(9)}$ :

There exists two constants $(\hat{P}_{44})^{(9)}$ and $(\hat{Q}_{44})^{(9)}$ which together with $(\hat{M}_{44})^{(9)}, (\hat{k}_{44})^{(9)}, (\hat{A}_{44})^{(9)} and\ (\hat{B}_{44})^{(9)}$ and the constants $(a_i)^{(9)}, (a_i^{'})^{(9)}, (b_i)^{(9)}, (b_i^{'})^{(9)}, (p_i)^{(9)}, (r_i)^{(9)}, i = 44,45,46,$ satisfy the inequalities

$$\frac{1}{(\hat{M}_{44})^{(9)}}[(a_i)^{(9)} + (a_i^{'})^{(9)} + (\hat{A}_{44})^{(9)} + (\hat{P}_{44})^{(9)}(\hat{k}_{44})^{(9)}] < 1$$

$$\frac{1}{(\hat{M}_{44})^{(9)}}[(b_i)^{(9)} + (b_i^{'})^{(9)} + (\hat{B}_{44})^{(9)} + (\hat{Q}_{44})^{(9)}(\hat{k}_{44})^{(9)}] < 1$$

**Theorem 1:** if the conditions above are fulfilled, there exists a solution satisfying the conditions 147

Definition of $G_i(0), T_i(0)$ :

$$G_i(t) \leq (\hat{P}_{13})^{(1)}e^{(\hat{M}_{13})^{(1)}t} \ , \ \boxed{G_i(0) = G_i^0 > 0}$$

$$T_i(t) \leq (\hat{Q}_{13})^{(1)}e^{(\hat{M}_{13})^{(1)}t} \ , \ \boxed{T_i(0) = T_i^0 > 0}$$

**Theorem 2 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions 148

**Definition of** $G_i(0), T_i(0)$

$$G_i(t) \leq (\hat{P}_{16})^{(2)}e^{(\hat{M}_{16})^{(2)}t} \ , \ G_i(0) = G_i^0 > 0$$

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| $T_i(t) \leq (\hat{Q}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ , $T_i(0) = T_i^0 > 0$ | |
| **Theorem 3 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>$G_i(t) \leq (\hat{P}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ , $G_i(0) = G_i^0 > 0$<br>$T_i(t) \leq (\hat{Q}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} t}$ , $T_i(0) = T_i^0 > 0$ | 149 |
| **Theorem 4 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \leq (\hat{P}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \leq (\hat{Q}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 150 |
| **Theorem 5 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \leq (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \leq (\hat{Q}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 151 |
| **Theorem 6 :** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \leq (\hat{P}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \leq (\hat{Q}_{32})^{(6)} e^{(\hat{M}_{32})^{(6)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 152 |
| **Theorem 7:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \leq (\hat{P}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br>$T_i(t) \leq (\hat{Q}_{36})^{(7)} e^{(\hat{M}_{36})^{(7)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 |
| **Theorem 8:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br>**Definition of** $G_i(0), T_i(0)$ :<br>$G_i(t) \leq (\hat{P}_{40})^{(8)} e^{(\hat{M}_{40})^{(8)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$ | 153 A |

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| $T_i(t) \leq (\hat{Q}_{40})^{(8)} e^{(\hat{M}_{40})^{(8)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | |
| **Theorem 9:** if the conditions above are fulfilled, there exists a solution satisfying the conditions<br><br>**Definition of** $G_i(0), T_i(0)$ :<br><br>$G_i(t) \leq (\hat{P}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} t}$ , $\boxed{G_i(0) = G_i^0 > 0}$<br><br>$T_i(t) \leq (\hat{Q}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} t}$ , $\boxed{T_i(0) = T_i^0 > 0}$ | 153 B |
| **Proof:** Consider operator $\mathcal{A}^{(1)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 154 |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{13})^{(1)}$ , $T_i^0 \leq (\hat{Q}_{13})^{(1)}$, | 155 |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ | 156 |
| $0 \leq T_i(t) - T_i^0 \leq (\hat{Q}_{13})^{(1)} e^{(\hat{M}_{13})^{(1)} t}$ | 157 |
| By<br><br>$\bar{G}_{13}(t) = G_{13}^0 + \int_0^t \left[ (a_{13})^{(1)} G_{14}(s_{(13)}) - \left( (a_{13}')^{(1)} + a_{13}'')^{(1)} (T_{14}(s_{(13)}), s_{(13)}) \right) G_{13}(s_{(13)}) \right] ds_{(13)}$ | 158 |
| $\bar{G}_{14}(t) = G_{14}^0 + \int_0^t \left[ (a_{14})^{(1)} G_{13}(s_{(13)}) - \left( (a_{14}')^{(1)} + (a_{14}'')^{(1)} (T_{14}(s_{(13)}), s_{(13)}) \right) G_{14}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{G}_{15}(t) = G_{15}^0 + \int_0^t \left[ (a_{15})^{(1)} G_{14}(s_{(13)}) - \left( (a_{15}')^{(1)} + (a_{15}'')^{(1)} (T_{14}(s_{(13)}), s_{(13)}) \right) G_{15}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{T}_{13}(t) = T_{13}^0 + \int_0^t \left[ (b_{13})^{(1)} T_{14}(s_{(13)}) - \left( (b_{13}')^{(1)} - (b_{13}'')^{(1)} (G(s_{(13)}), s_{(13)}) \right) T_{13}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{T}_{14}(t) = T_{14}^0 + \int_0^t \left[ (b_{14})^{(1)} T_{13}(s_{(13)}) - \left( (b_{14}')^{(1)} - (b_{14}'')^{(1)} (G(s_{(13)}), s_{(13)}) \right) T_{14}(s_{(13)}) \right] ds_{(13)}$ | |
| $\bar{T}_{15}(t) = T_{15}^0 + \int_0^t \left[ (b_{15})^{(1)} T_{14}(s_{(13)}) - \left( (b_{15}')^{(1)} - (b_{15}'')^{(1)} (G(s_{(13)}), s_{(13)}) \right) T_{15}(s_{(13)}) \right] ds_{(13)}$<br><br>Where $s_{(13)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br><br>Consider operator $\mathcal{A}^{(2)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 159 |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \leq (\hat{P}_{16})^{(2)}$ , $T_i^0 \leq (\hat{Q}_{16})^{(2)}$, | |
| $0 \leq G_i(t) - G_i^0 \leq (\hat{P}_{16})^{(2)} e^{(\hat{M}_{16})^{(2)} t}$ | |

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| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{16})^{(2)} e^{(\bar{M}_{16})^{(2)} t}$ | |
| By<br>$\bar{G}_{16}(t) = G_{16}^0 + \int_0^t \left[ (a_{16})^{(2)} G_{17}(s_{(16)}) - \left( (a_{16}')^{(2)} + a_{16}'')^{(2)} (T_{17}(s_{(16)}), s_{(16)}) \right) G_{16}(s_{(16)}) \right] ds_{(16)}$ | 160 |
| $\bar{G}_{17}(t) = G_{17}^0 + \int_0^t \left[ (a_{17})^{(2)} G_{16}(s_{(16)}) - \left( (a_{17}')^{(2)} + (a_{17}'')^{(2)} (T_{17}(s_{(16)}), s_{(17)}) \right) G_{17}(s_{(16)}) \right] ds_{(16)}$ | |
| $\bar{G}_{18}(t) = G_{18}^0 + \int_0^t \left[ (a_{18})^{(2)} G_{17}(s_{(16)}) - \left( (a_{18}')^{(2)} + (a_{18}'')^{(2)} (T_{17}(s_{(16)}), s_{(16)}) \right) G_{18}(s_{(16)}) \right] ds_{(16)}$ | |
| $\bar{T}_{16}(t) = T_{16}^0 + \int_0^t \left[ (b_{16})^{(2)} T_{17}(s_{(16)}) - \left( (b_{16}')^{(2)} - (b_{16}'')^{(2)} (G_{19}(s_{(16)}), s_{(16)}) \right) T_{16}(s_{(16)}) \right] ds_{(16)}$ | |
| $\bar{T}_{17}(t) = T_{17}^0 + \int_0^t \left[ (b_{17})^{(2)} T_{16}(s_{(16)}) - \left( (b_{17}')^{(2)} - (b_{17}'')^{(2)} (G_{19}(s_{(16)}), s_{(16)}) \right) T_{17}(s_{(16)}) \right] ds_{(16)}$ | |
| $\bar{T}_{18}(t) = T_{18}^0 + \int_0^t \left[ (b_{18})^{(2)} T_{17}(s_{(16)}) - \left( (b_{18}')^{(2)} - (b_{18}'')^{(2)} (G_{19}(s_{(16)}), s_{(16)}) \right) T_{18}(s_{(16)}) \right] ds_{(16)}$<br>Where $s_{(16)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(3)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{20})^{(3)}$, $T_i^0 \le (\hat{Q}_{20})^{(3)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{20})^{(3)} e^{(\bar{M}_{20})^{(3)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{20})^{(3)} e^{(\bar{M}_{20})^{(3)} t}$ | |
| By<br>$\bar{G}_{20}(t) = G_{20}^0 + \int_0^t \left[ (a_{20})^{(3)} G_{21}(s_{(20)}) - \left( (a_{20}')^{(3)} + a_{20}'')^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{20}(s_{(20)}) \right] ds_{(20)}$ | 161 |
| $\bar{G}_{21}(t) = G_{21}^0 + \int_0^t \left[ (a_{21})^{(3)} G_{20}(s_{(20)}) - \left( (a_{21}')^{(3)} + (a_{21}'')^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{21}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{G}_{22}(t) = G_{22}^0 + \int_0^t \left[ (a_{22})^{(3)} G_{21}(s_{(20)}) - \left( (a_{22}')^{(3)} + (a_{22}'')^{(3)} (T_{21}(s_{(20)}), s_{(20)}) \right) G_{22}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{20}(t) = T_{20}^0 + \int_0^t \left[ (b_{20})^{(3)} T_{21}(s_{(20)}) - \left( (b_{20}')^{(3)} - (b_{20}'')^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{20}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{21}(t) = T_{21}^0 + \int_0^t \left[ (b_{21})^{(3)} T_{20}(s_{(20)}) - \left( (b_{21}')^{(3)} - (b_{21}'')^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{21}(s_{(20)}) \right] ds_{(20)}$ | |
| $\bar{T}_{22}(t) = T_{22}^0 + \int_0^t \left[ (b_{22})^{(3)} T_{21}(s_{(20)}) - \left( (b_{22}')^{(3)} - (b_{22}'')^{(3)} (G_{23}(s_{(20)}), s_{(20)}) \right) T_{22}(s_{(20)}) \right] ds_{(20)}$<br>Where $s_{(20)}$ is the integrand that is integrated over an interval $(0, t)$ | |

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| **Proof:** Consider operator $\mathcal{A}^{(4)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i$: $\mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \le (\hat{P}_{24})^{(4)}$ , $T_i^0 \le (\hat{Q}_{24})^{(4)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} t}$ | |
| By<br>$\bar{G}_{24}(t) = G_{24}^0 + \int_0^t \left[ (a_{24})^{(4)} G_{25}(s_{(24)}) - \left( (a'_{24})^{(4)} + a''_{24})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{24}(s_{(24)}) \right] ds_{(24)}$ | 162 |
| $\bar{G}_{25}(t) = G_{25}^0 + \int_0^t \left[ (a_{25})^{(4)} G_{24}(s_{(24)}) - \left( (a'_{25})^{(4)} + (a''_{25})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{25}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{G}_{26}(t) = G_{26}^0 + \int_0^t \left[ (a_{26})^{(4)} G_{25}(s_{(24)}) - \left( (a'_{26})^{(4)} + (a''_{26})^{(4)} (T_{25}(s_{(24)}), s_{(24)}) \right) G_{26}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{24}(t) = T_{24}^0 + \int_0^t \left[ (b_{24})^{(4)} T_{25}(s_{(24)}) - \left( (b'_{24})^{(4)} - (b''_{24})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{24}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{25}(t) = T_{25}^0 + \int_0^t \left[ (b_{25})^{(4)} T_{24}(s_{(24)}) - \left( (b'_{25})^{(4)} - (b''_{25})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{25}(s_{(24)}) \right] ds_{(24)}$ | |
| $\bar{T}_{26}(t) = T_{26}^0 + \int_0^t \left[ (b_{26})^{(4)} T_{25}(s_{(24)}) - \left( (b'_{26})^{(4)} - (b''_{26})^{(4)} (G_{27}(s_{(24)}), s_{(24)}) \right) T_{26}(s_{(24)}) \right] ds_{(24)}$<br>Where $s_{(24)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:** Consider operator $\mathcal{A}^{(5)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i$: $\mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \le (\hat{P}_{28})^{(5)}$ , $T_i^0 \le (\hat{Q}_{28})^{(5)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} t}$ | |
| By<br>$\bar{G}_{28}(t) = G_{28}^0 + \int_0^t \left[ (a_{28})^{(5)} G_{29}(s_{(28)}) - \left( (a'_{28})^{(5)} + a''_{28})^{(5)} (T_{29}(s_{(28)}), s_{(28)}) \right) G_{28}(s_{(28)}) \right] ds_{(28)}$ | 163 |
| $\bar{G}_{29}(t) = G_{29}^0 + \int_0^t \left[ (a_{29})^{(5)} G_{28}(s_{(28)}) - \left( (a'_{29})^{(5)} + (a''_{29})^{(5)} (T_{29}(s_{(28)}), s_{(28)}) \right) G_{29}(s_{(28)}) \right] ds_{(28)}$ | |
| $\bar{G}_{30}(t) = G_{30}^0 + \int_0^t \left[ (a_{30})^{(5)} G_{29}(s_{(28)}) - \left( (a'_{30})^{(5)} + (a''_{30})^{(5)} (T_{29}(s_{(28)}), s_{(28)}) \right) G_{30}(s_{(28)}) \right] ds_{(28)}$ | |
| $\bar{T}_{28}(t) = T_{28}^0 + \int_0^t \left[ (b_{28})^{(5)} T_{29}(s_{(28)}) - \left( (b'_{28})^{(5)} - (b''_{28})^{(5)} (G_{31}(s_{(28)}), s_{(28)}) \right) T_{28}(s_{(28)}) \right] ds_{(28)}$ | |

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| $\bar{T}_{29}(t) = T_{29}^0 + \int_0^t \left[(b_{29})^{(5)} T_{28}(s_{(28)}) - \left((b'_{29})^{(5)} - (b''_{29})^{(5)}(G_{31}(s_{(28)}), s_{(28)})\right) T_{29}(s_{(28)})\right] ds_{(28)}$ | |
| $\bar{\mathrm{T}}_{30}(\mathrm{t}) = \mathrm{T}_{30}^0 + \int_0^t \left[(b_{30})^{(5)} T_{29}(s_{(28)}) - \left((b'_{30})^{(5)} - (b''_{30})^{(5)}(G_{31}(s_{(28)}), s_{(28)})\right) T_{30}(s_{(28)})\right] ds_{(28)}$<br>Where $s_{(28)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(6)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \le (\hat{P}_{32})^{(6)}$ , $T_i^0 \le (\hat{Q}_{32})^{(6)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{32})^{(6)} e^{(\bar{M}_{32})^{(6)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{32})^{(6)} e^{(\bar{M}_{32})^{(6)} t}$ | |
| By<br>$\bar{G}_{32}(t) = G_{32}^0 + \int_0^t \left[(a_{32})^{(6)} G_{33}(s_{(32)}) - \left((a'_{32})^{(6)} + a''_{32})^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{32}(s_{(32)})\right] ds_{(32)}$ | 164 |
| $\bar{G}_{33}(t) = G_{33}^0 + \int_0^t \left[(a_{33})^{(6)} G_{32}(s_{(32)}) - \left((a'_{33})^{(6)} + (a''_{33})^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{33}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{G}_{34}(t) = G_{34}^0 + \int_0^t \left[(a_{34})^{(6)} G_{33}(s_{(32)}) - \left((a'_{34})^{(6)} + (a''_{34})^{(6)}(T_{33}(s_{(32)}), s_{(32)})\right) G_{34}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{T}_{32}(t) = T_{32}^0 + \int_0^t \left[(b_{32})^{(6)} T_{33}(s_{(32)}) - \left((b'_{32})^{(6)} - (b''_{32})^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{32}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{T}_{33}(t) = T_{33}^0 + \int_0^t \left[(b_{33})^{(6)} T_{32}(s_{(32)}) - \left((b'_{33})^{(6)} - (b''_{33})^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{33}(s_{(32)})\right] ds_{(32)}$ | |
| $\bar{\mathrm{T}}_{34}(\mathrm{t}) = \mathrm{T}_{34}^0 + \int_0^t \left[(b_{34})^{(6)} T_{33}(s_{(32)}) - \left((b'_{34})^{(6)} - (b''_{34})^{(6)}(G_{35}(s_{(32)}), s_{(32)})\right) T_{34}(s_{(32)})\right] ds_{(32)}$<br>Where $s_{(32)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(7)}$ defined on the space of sextuples of continuous functions $G_i$ , $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$ , $T_i(0) = T_i^0$ , $G_i^0 \le (\hat{P}_{36})^{(7)}$ , $T_i^0 \le (\hat{Q}_{36})^{(7)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{36})^{(7)} e^{(\bar{M}_{36})^{(7)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{36})^{(7)} e^{(\bar{M}_{36})^{(7)} t}$ | |

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| By<br>$\bar{G}_{36}(t) = G_{36}^{0} + \int_0^t \left[(a_{36})^{(7)} G_{37}(s_{(36)}) - \left((a_{36}^{'})^{(7)} + a_{36}^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{36}(s_{(36)})\right] ds_{(36)}$ | 165 |
| $\bar{G}_{37}(t) = G_{37}^{0} + \int_0^t \left[(a_{37})^{(7)} G_{36}(s_{(36)}) - \left((a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{37}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{G}_{38}(t) = G_{38}^{0} + \int_0^t \left[(a_{38})^{(7)} G_{37}(s_{(36)}) - \left((a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37}(s_{(36)}), s_{(36)})\right) G_{38}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{36}(t) = T_{36}^{0} + \int_0^t \left[(b_{36})^{(7)} T_{37}(s_{(36)}) - \left((b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{36}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{37}(t) = T_{37}^{0} + \int_0^t \left[(b_{37})^{(7)} T_{36}(s_{(36)}) - \left((b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{37}(s_{(36)})\right] ds_{(36)}$ | |
| $\bar{T}_{38}(t) = T_{38}^{0} + \int_0^t \left[(b_{38})^{(7)} T_{37}(s_{(36)}) - \left((b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}(G_{39}(s_{(36)}), s_{(36)})\right) T_{38}(s_{(36)})\right] ds_{(36)}$<br>Where $s_{(36)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(8)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{40})^{(8)}$, $T_i^0 \le (\hat{Q}_{40})^{(8)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{40})^{(8)} e^{(\bar{M}_{40})^{(8)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{40})^{(8)} e^{(\bar{M}_{40})^{(8)} t}$ | |
| By<br>$\bar{G}_{40}(t) = G_{40}^{0} + \int_0^t \left[(a_{40})^{(8)} G_{41}(s_{(40)}) - \left((a_{40}^{'})^{(8)} + a_{40}^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\right) G_{40}(s_{(40)})\right] ds_{(40)}$ | 166 |
| $\bar{G}_{41}(t) = G_{41}^{0} + \int_0^t \left[(a_{41})^{(8)} G_{40}(s_{(40)}) - \left((a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\right) G_{41}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{G}_{42}(t) = G_{42}^{0} + \int_0^t \left[(a_{42})^{(8)} G_{41}(s_{(40)}) - \left((a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41}(s_{(40)}), s_{(40)})\right) G_{42}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{T}_{40}(t) = T_{40}^{0} + \int_0^t \left[(b_{40})^{(8)} T_{41}(s_{(40)}) - \left((b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}(G_{43}(s_{(40)}), s_{(40)})\right) T_{40}(s_{(40)})\right] ds_{(40)}$ | |
| $\bar{T}_{41}(t) = T_{41}^{0} + \int_0^t \left[(b_{41})^{(8)} T_{40}(s_{(40)}) - \left((b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}(G_{43}(s_{(40)}), s_{(40)})\right) T_{41}(s_{(40)})\right] ds_{(40)}$ | |

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| $\overline{T}_{42}(t) = T_{42}^0 + \int_0^t \left[(b_{42})^{(8)} T_{41}(s_{(40)}) - \left((b'_{42})^{(8)} - (b''_{42})^{(8)}(G_{43}(s_{(40)}), s_{(40)})\right) T_{42}(s_{(40)})\right] ds_{(40)}$<br><br>Where $s_{(40)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| **Proof:**<br>Consider operator $\mathcal{A}^{(9)}$ defined on the space of sextuples of continuous functions $G_i$, $T_i: \mathbb{R}_+ \to \mathbb{R}_+$ which satisfy | 166<br>A |
| $G_i(0) = G_i^0$, $T_i(0) = T_i^0$, $G_i^0 \le (\hat{P}_{44})^{(9)}$, $T_i^0 \le (\hat{Q}_{44})^{(9)}$, | |
| $0 \le G_i(t) - G_i^0 \le (\hat{P}_{44})^{(9)} e^{(\bar{M}_{44})^{(9)} t}$ | |
| $0 \le T_i(t) - T_i^0 \le (\hat{Q}_{44})^{(9)} e^{(\bar{M}_{44})^{(9)} t}$ | |
| By<br><br>$\bar{G}_{44}(t) = G_{44}^0 + \int_0^t \left[(a_{44})^{(9)} G_{45}(s_{(44)}) - \left((a'_{44})^{(9)} + a''_{44})^{(9)}(T_{45}(s_{(44)}), s_{(44)})\right) G_{44}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{G}_{45}(t) = G_{45}^0 + \int_0^t \left[(a_{45})^{(9)} G_{44}(s_{(44)}) - \left((a'_{45})^{(9)} + (a''_{45})^{(9)}(T_{45}(s_{(44)}), s_{(44)})\right) G_{45}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{G}_{46}(t) = G_{46}^0 + \int_0^t \left[(a_{46})^{(9)} G_{45}(s_{(44)}) - \left((a'_{46})^{(9)} + (a''_{46})^{(9)}(T_{45}(s_{(44)}), s_{(44)})\right) G_{46}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{44}(t) = T_{44}^0 + \int_0^t \left[(b_{44})^{(9)} T_{45}(s_{(44)}) - \left((b'_{44})^{(9)} - (b''_{44})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{44}(s_{(44)})\right] ds_{(44)}$ | |
| $\bar{T}_{45}(t) = T_{45}^0 + \int_0^t \left[(b_{45})^{(9)} T_{44}(s_{(44)}) - \left((b'_{45})^{(9)} - (b''_{45})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{45}(s_{(44)})\right] ds_{(44)}$ | |
| $\overline{T}_{46}(t) = T_{46}^0 + \int_0^t \left[(b_{46})^{(9)} T_{45}(s_{(44)}) - \left((b'_{46})^{(9)} - (b''_{46})^{(9)}(G_{47}(s_{(44)}), s_{(44)})\right) T_{46}(s_{(44)})\right] ds_{(44)}$<br><br>Where $s_{(44)}$ is the integrand that is integrated over an interval $(0, t)$ | |
| The operator $\mathcal{A}^{(1)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{13}(t) \le G_{13}^0 + \int_0^t \left[(a_{13})^{(1)} \left(G_{14}^0 + (\hat{P}_{13})^{(1)} e^{(\bar{M}_{13})^{(1)} s_{(13)}}\right)\right] ds_{(13)} =$<br><br>$\left(1 + (a_{13})^{(1)} t\right) G_{14}^0 + \frac{(a_{13})^{(1)} (\hat{P}_{13})^{(1)}}{(\bar{M}_{13})^{(1)}} \left(e^{(\bar{M}_{13})^{(1)} t} - 1\right)$ | 167 |
| From which it follows that<br><br>$(G_{13}(t) - G_{13}^0) e^{-(\bar{M}_{13})^{(1)} t} \le \frac{(a_{13})^{(1)}}{(\bar{M}_{13})^{(1)}} \left[\left((\hat{P}_{13})^{(1)} + G_{14}^0\right) e^{\left(-\frac{(\hat{P}_{13})^{(1)} + G_{14}^0}{G_{14}^0}\right)} + (\hat{P}_{13})^{(1)}\right]$<br><br>$(G_i^0)$ is as defined in the statement of theorem 1 | 168 |
| Analogous inequalities hold also for $G_{14}, G_{15}, T_{13}, T_{14}, T_{15}$ | |

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| The operator $\mathcal{A}^{(2)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that | |
| $G_{16}(t) \le G_{16}^0 + \int_0^t \left[(a_{16})^{(2)}\left(G_{17}^0 + (\hat{P}_{16})^{(6)} e^{(\hat{M}_{16})^{(2)} s_{(16)}}\right)\right] ds_{(16)} =$ $\left(1 + (a_{16})^{(2)} t\right) G_{17}^0 + \frac{(a_{16})^{(2)} (\hat{P}_{16})^{(2)}}{(\hat{M}_{16})^{(2)}}\left(e^{(\hat{M}_{16})^{(2)} t} - 1\right)$ | 169 |
| From which it follows that $(G_{16}(t) - G_{16}^0) e^{-(\hat{M}_{16})^{(2)} t} \le \frac{(a_{16})^{(2)}}{(\hat{M}_{16})^{(2)}}\left[\left((\hat{P}_{16})^{(2)} + G_{17}^0\right) e^{\left(-\frac{(\hat{P}_{16})^{(2)} + G_{17}^0}{G_{17}^0}\right)} + (\hat{P}_{16})^{(2)}\right]$ | 170 |
| Analogous inequalities hold also for $G_{17}, G_{18}, T_{16}, T_{17}, T_{18}$ | |
| The operator $\mathcal{A}^{(3)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{20}(t) \le G_{20}^0 + \int_0^t \left[(a_{20})^{(3)}\left(G_{21}^0 + (\hat{P}_{20})^{(3)} e^{(\hat{M}_{20})^{(3)} s_{(20)}}\right)\right] ds_{(20)} =$ $\left(1 + (a_{20})^{(3)} t\right) G_{21}^0 + \frac{(a_{20})^{(3)} (\hat{P}_{20})^{(3)}}{(\hat{M}_{20})^{(3)}}\left(e^{(\hat{M}_{20})^{(3)} t} - 1\right)$ | 171 |
| From which it follows that $(G_{20}(t) - G_{20}^0) e^{-(\hat{M}_{20})^{(3)} t} \le \frac{(a_{20})^{(3)}}{(\hat{M}_{20})^{(3)}}\left[\left((\hat{P}_{20})^{(3)} + G_{21}^0\right) e^{\left(-\frac{(\hat{P}_{20})^{(3)} + G_{21}^0}{G_{21}^0}\right)} + (\hat{P}_{20})^{(3)}\right]$ | 172 |
| Analogous inequalities hold also for $G_{21}, G_{22}, T_{20}, T_{21}, T_{22}$ | |
| The operator $\mathcal{A}^{(4)}$ maps the space of functions satisfying into itself .Indeed it is obvious that $G_{24}(t) \le G_{24}^0 + \int_0^t \left[(a_{24})^{(4)}\left(G_{25}^0 + (\hat{P}_{24})^{(4)} e^{(\hat{M}_{24})^{(4)} s_{(24)}}\right)\right] ds_{(24)} =$ $\left(1 + (a_{24})^{(4)} t\right) G_{25}^0 + \frac{(a_{24})^{(4)} (\hat{P}_{24})^{(4)}}{(\hat{M}_{24})^{(4)}}\left(e^{(\hat{M}_{24})^{(4)} t} - 1\right)$ | 173 |
| From which it follows that $(G_{24}(t) - G_{24}^0) e^{-(\hat{M}_{24})^{(4)} t} \le \frac{(a_{24})^{(4)}}{(\hat{M}_{24})^{(4)}}\left[\left((\hat{P}_{24})^{(4)} + G_{25}^0\right) e^{\left(-\frac{(\hat{P}_{24})^{(4)} + G_{25}^0}{G_{25}^0}\right)} + (\hat{P}_{24})^{(4)}\right]$ $(G_i^0)$ is as defined in the statement of theorem 4 | 174 |
| The operator $\mathcal{A}^{(5)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that $G_{28}(t) \le G_{28}^0 + \int_0^t \left[(a_{28})^{(5)}\left(G_{29}^0 + (\hat{P}_{28})^{(5)} e^{(\hat{M}_{28})^{(5)} s_{(28)}}\right)\right] ds_{(28)} =$ $\left(1 + (a_{28})^{(5)} t\right) G_{29}^0 + \frac{(a_{28})^{(5)} (\hat{P}_{28})^{(5)}}{(\hat{M}_{28})^{(5)}}\left(e^{(\hat{M}_{28})^{(5)} t} - 1\right)$ | |

From which it follows that

$$(G_{28}(t)-G_{28}^{0})e^{-(\bar{M}_{28})^{(5)}t} \le \frac{(a_{28})^{(5)}}{(\bar{M}_{28})^{(5)}}\left[((\hat{P}_{28})^{(5)}+G_{29}^{0})e^{\left(-\frac{(\hat{P}_{28})^{(5)}+G_{29}^{0}}{G_{29}^{0}}\right)}+(\hat{P}_{28})^{(5)}\right]$$

$(G_i^0)$ is as defined in the statement of theorem 5 175

The operator $\mathcal{A}^{(6)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that 176

$$G_{32}(t) \le G_{32}^{0}+\int_0^t\left[(a_{32})^{(6)}\left(G_{33}^{0}+(\hat{P}_{32})^{(6)}e^{(\bar{M}_{32})^{(6)}s_{(32)}}\right)\right]ds_{(32)} =$$

$$\left(1+(a_{32})^{(6)}t\right)G_{33}^{0}+\frac{(a_{32})^{(6)}(\hat{P}_{32})^{(6)}}{(\bar{M}_{32})^{(6)}}\left(e^{(\bar{M}_{32})^{(6)}t}-1\right)$$

From which it follows that 177

$$(G_{32}(t)-G_{32}^{0})e^{-(\bar{M}_{32})^{(6)}t} \le \frac{(a_{32})^{(6)}}{(\bar{M}_{32})^{(6)}}\left[((\hat{P}_{32})^{(6)}+G_{33}^{0})e^{\left(-\frac{(\hat{P}_{32})^{(6)}+G_{33}^{0}}{G_{33}^{0}}\right)}+(\hat{P}_{32})^{(6)}\right]$$

$(G_i^0)$ is as defined in the statement of theorem 6

Analogous inequalities hold also for $G_{25}, G_{26}, T_{24}, T_{25}, T_{26}$

(h) The operator $\mathcal{A}^{(7)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that 178

$$G_{36}(t) \le G_{36}^{0}+\int_0^t\left[(a_{36})^{(7)}\left(G_{37}^{0}+(\hat{P}_{36})^{(7)}e^{(\bar{M}_{36})^{(7)}s_{(36)}}\right)\right]ds_{(36)} =$$

$$\left(1+(a_{36})^{(7)}t\right)G_{37}^{0}+\frac{(a_{36})^{(7)}(\hat{P}_{36})^{(7)}}{(\bar{M}_{36})^{(7)}}\left(e^{(\bar{M}_{36})^{(7)}t}-1\right)$$

From which it follows that

$$(G_{36}(t)-G_{36}^{0})e^{-(\bar{M}_{36})^{(7)}t} \le \frac{(a_{36})^{(7)}}{(\bar{M}_{36})^{(7)}}\left[((\hat{P}_{36})^{(7)}+G_{37}^{0})e^{\left(-\frac{(\hat{P}_{36})^{(7)}+G_{37}^{0}}{G_{37}^{0}}\right)}+(\hat{P}_{36})^{(7)}\right]$$

$(G_i^0)$ is as defined in the statement of theorem 7

The operator $\mathcal{A}^{(8)}$ maps the space of functions satisfying Equations into itself .Indeed it is obvious that

$$G_{40}(t) \le G_{40}^{0}+\int_0^t\left[(a_{40})^{(8)}\left(G_{41}^{0}+(\hat{P}_{40})^{(8)}e^{(\bar{M}_{40})^{(8)}s_{(40)}}\right)\right]ds_{(40)} =$$

$$\left(1+(a_{40})^{(8)}t\right)G_{41}^{0}+\frac{(a_{40})^{(8)}(\hat{P}_{40})^{(8)}}{(\bar{M}_{40})^{(8)}}\left(e^{(\bar{M}_{40})^{(8)}t}-1\right)$$ 180

From which it follows that

$$(G_{40}(t)-G_{40}^{0})e^{-(\bar{M}_{40})^{(8)}t} \le \frac{(a_{40})^{(8)}}{(\bar{M}_{40})^{(8)}}\left[((\hat{P}_{40})^{(8)}+G_{41}^{0})e^{\left(-\frac{(\hat{P}_{40})^{(8)}+G_{41}^{0}}{G_{41}^{0}}\right)}+(\hat{P}_{40})^{(8)}\right]$$ 181

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| $(G_i^0)$ is as defined in the statement of theorem 8<br>Analogous inequalities hold also for $G_{41}, G_{42}, T_{40}, T_{41}, T_{42}$ | |
| The operator $\mathcal{A}^{(9)}$ maps the space of functions satisfying 34,35,36 into itself .Indeed it is obvious that<br>$G_{44}(t) \le G_{44}^0 + \int_0^t \left[(a_{44})^{(9)}\left(G_{45}^0 + (\hat{P}_{44})^{(9)} e^{(\hat{M}_{44})^{(9)} s_{(44)}}\right)\right] ds_{(44)} =$<br>$\left(1 + (a_{44})^{(9)} t\right) G_{45}^0 + \frac{(a_{44})^{(9)} (\hat{P}_{44})^{(9)}}{(\hat{M}_{44})^{(9)}}\left(e^{(\hat{M}_{44})^{(9)} t} - 1\right)$ | |
| From which it follows that<br>$(G_{44}(t) - G_{44}^0) e^{-(\hat{M}_{44})^{(9)} t} \le \frac{(a_{44})^{(9)}}{(\hat{M}_{44})^{(9)}}\left[\left((\hat{P}_{44})^{(9)} + G_{45}^0\right) e^{\left(-\frac{(\hat{P}_{44})^{(9)} + G_{45}^0}{G_{45}^0}\right)} + (\hat{P}_{44})^{(9)}\right]$<br>$(G_i^0)$ is as defined in the statement of theorem 9<br>Analogous inequalities hold also for $G_{45}, G_{46}, T_{44}, T_{45}, T_{46}$ | |
| It is now sufficient to take $\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}, \frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}} < 1$ and to choose<br>$(\hat{P}_{13})^{(1)}$ and $(\hat{Q}_{13})^{(1)}$ large to have | 182 |
| $\frac{(a_i)^{(1)}}{(\hat{M}_{13})^{(1)}}\left[(\hat{P}_{13})^{(1)} + \left((\hat{P}_{13})^{(1)} + G_j^0\right) e^{-\left(\frac{(\hat{P}_{13})^{(1)} + G_j^0}{G_j^0}\right)}\right] \le (\hat{P}_{13})^{(1)}$ | 183 |
| $\frac{(b_i)^{(1)}}{(\hat{M}_{13})^{(1)}}\left[\left((\hat{Q}_{13})^{(1)} + T_j^0\right) e^{-\left(\frac{(\hat{Q}_{13})^{(1)} + T_j^0}{T_j^0}\right)} + (\hat{Q}_{13})^{(1)}\right] \le (\hat{Q}_{13})^{(1)}$ | 184 |
| In order that the operator $\mathcal{A}^{(1)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(1)}$ is a contraction with respect to the metric<br>$d\left(\left(G^{(1)}, T^{(1)}\right), \left(G^{(2)}, T^{(2)}\right)\right) =$<br>$\sup_i \{\max_{t \in \mathbb{R}_+} \left|G_i^{(1)}(t) - G_i^{(2)}(t)\right| e^{-(\hat{M}_{13})^{(1)} t}, \max_{t \in \mathbb{R}_+} \left|T_i^{(1)}(t) - T_i^{(2)}(t)\right| e^{-(\hat{M}_{13})^{(1)} t}\}$ | 185 |
| Indeed if we denote<br>**Definition of** $\tilde{G}, \tilde{T}$ : $\left(\tilde{G}, \tilde{T}\right) = \mathcal{A}^{(1)}(G, T)$<br>It results | |

$$\left|\tilde{G}_{13}^{(1)}-\tilde{G}_{i}^{(2)}\right| \leq \int_0^t (a_{13})^{(1)} \left|G_{14}^{(1)}-G_{14}^{(2)}\right| e^{-(\widehat{M}_{13})^{(1)}s_{(13)}} e^{(\widehat{M}_{13})^{(1)}s_{(13)}}\, ds_{(13)} +$$

$$\int_0^t \{(a_{13}')^{(1)} \left|G_{13}^{(1)}-G_{13}^{(2)}\right| e^{-(\widehat{M}_{13})^{(1)}s_{(13)}} e^{-(\widehat{M}_{13})^{(1)}s_{(13)}} +$$

$$(a_{13}'')^{(1)}\left(T_{14}^{(1)}, s_{(13)}\right)\left|G_{13}^{(1)}-G_{13}^{(2)}\right| e^{-(\widehat{M}_{13})^{(1)}s_{(13)}} e^{(\widehat{M}_{13})^{(1)}s_{(13)}} +$$

$$G_{13}^{(2)}\left|(a_{13}'')^{(1)}\left(T_{14}^{(1)}, s_{(13)}\right)-(a_{13}'')^{(1)}\left(T_{14}^{(2)}, s_{(13)}\right)\right|\ e^{-(\widehat{M}_{13})^{(1)}s_{(13)}} e^{(\widehat{M}_{13})^{(1)}s_{(13)}}\} ds_{(13)}$$

Where $s_{(13)}$ represents integrand that is integrated over the interval $[0, t]$

From the hypotheses it follows

$$\left|G^{(1)}-G^{(2)}\right| e^{-(\widehat{M}_{13})^{(1)}t} \leq$$
$$\frac{1}{(\widehat{M}_{13})^{(1)}}\left((a_{13})^{(1)} + (a_{13}')^{(1)} + (\widehat{A}_{13})^{(1)} + (\widehat{P}_{13})^{(1)}(\widehat{k}_{13})^{(1)}\right) d\left(\left(G^{(1)}, T^{(1)};\ G^{(2)}, T^{(2)}\right)\right)$$

And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows

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**Remark 1:** The fact that we supposed $(a_{13}'')^{(1)}$ and $(b_{13}'')^{(1)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{13})^{(1)} e^{(\widehat{M}_{13})^{(1)}t}$ *and* $(\widehat{Q}_{13})^{(1)} e^{(\widehat{M}_{13})^{(1)}t}$ respectively of $\mathbb{R}_+$.

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i'')^{(1)}$ and $(b_i'')^{(1)}, i = 13,14,15$ depend only on $T_{14}$ and respectively on $G(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition.

**Remark 2:** There does not exist any $t$ where $G_i(t) = 0$ *and* $T_i(t) = 0$

From 19 to 24 it results

$$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i')^{(1)} - (a_i'')^{(1)}(T_{14}(s_{(13)}), s_{(13)})\} ds_{(13)}\right]} \geq 0$$

$$T_i(t) \geq T_i^0 e^{\left(-(b_i')^{(1)}t\right)} > 0 \quad \text{for t} > 0$$

**Definition of** $\left((\widehat{M}_{13})^{(1)}\right)_1, \left((\widehat{M}_{13})^{(1)}\right)_2$ *and* $\left((\widehat{M}_{13})^{(1)}\right)_3$ :

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**Remark 3**: if $G_{13}$ is bounded, the same property have also $G_{14}$ *and* $G_{15}$ . indeed if

$G_{13} < (\widehat{M}_{13})^{(1)}$ it follows $\frac{dG_{14}}{dt} \leq \left((\widehat{M}_{13})^{(1)}\right)_1 - (a_{14}')^{(1)} G_{14}$ and by integrating

$$G_{14} \leq \left((\widehat{M}_{13})^{(1)}\right)_2 = G_{14}^0 + 2(a_{14})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_1 / (a_{14}')^{(1)}$$

In the same way , one can obtain

$$G_{15} \leq \left((\widehat{M}_{13})^{(1)}\right)_3 = G_{15}^0 + 2(a_{15})^{(1)}\left((\widehat{M}_{13})^{(1)}\right)_2 / (a_{15}')^{(1)}$$

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| If $G_{14}$ *or* $G_{15}$ is bounded, the same property follows for $G_{13}$ , $G_{15}$ and $G_{13}$ , $G_{14}$ respectively. | |
| **Remark 4:** If $G_{13}$ *is* bounded, from below, the same property holds for $G_{14}$ *and* $G_{15}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{14}$ is bounded from below. | 188 |
| **Remark 5:** If $T_{13}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(1)}(G(t),t)) = (b_{14}^{'})^{(1)}$ then $T_{14} \to \infty$.<br><br>**Definition of** $(m)^{(1)}$ and $\varepsilon_1$ :<br><br>Indeed let $t_1$ be so that for $t > t_1$<br><br>$(b_{14})^{(1)} - (b_i^{''})^{(1)}(G(t),t) < \varepsilon_1, T_{13}(t) > (m)^{(1)}$ | 189 |
| Then $\frac{dT_{14}}{dt} \geq (a_{14})^{(1)}(m)^{(1)} - \varepsilon_1 T_{14}$ which leads to<br><br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{\varepsilon_1}\right)(1 - e^{-\varepsilon_1 t}) + T_{14}^0 e^{-\varepsilon_1 t}$ If we take t such that $e^{-\varepsilon_1 t} = \frac{1}{2}$ it results<br><br>$T_{14} \geq \left(\frac{(a_{14})^{(1)}(m)^{(1)}}{2}\right), \quad t = log\frac{2}{\varepsilon_1}$ By taking now $\varepsilon_1$ sufficiently small one sees that $T_{14}$ is unbounded.<br>The same property holds for $T_{15}$ if $\lim_{t\to\infty}(b_{15}^{''})^{(1)}(G(t),t) = (b_{15}^{'})^{(1)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | |
| It is now sufficient to take $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}$ , $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}} < 1$ and to choose<br><br>$(\hat{P}_{16})^{(2)}$ *and* $(\hat{Q}_{16})^{(2)}$ large to have | 190 |
| $\frac{(a_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[(\hat{P}_{16})^{(2)} + \left((\hat{P}_{16})^{(2)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{16})^{(2)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{16})^{(2)}$ | 191 |
| $\frac{(b_i)^{(2)}}{(\hat{M}_{16})^{(2)}}\left[\left((\hat{Q}_{16})^{(2)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{16})^{(2)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{16})^{(2)}\right] \leq (\hat{Q}_{16})^{(2)}$ | 192 |
| In order that the operator $\mathcal{A}^{(2)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 193 |
| The operator $\mathcal{A}^{(2)}$ is a contraction with respect to the metric<br><br>$d\left(((G_{19})^{(1)},(T_{19})^{(1)}),((G_{19})^{(2)},(T_{19})^{(2)})\right) =$<br><br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left|G_i^{(1)}(t) - G_i^{(2)}(t)\right|e^{-(\hat{M}_{16})^{(2)}t}, \max_{t\in\mathbb{R}_+}\left|T_i^{(1)}(t) - T_i^{(2)}(t)\right|e^{-(\hat{M}_{16})^{(2)}t}\}$ | 194 |
| Indeed if we denote | 195 |

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| **Definition of** $\widetilde{G_{19}}, \widetilde{T_{19}}$ : $\left(\widetilde{G_{19}}, \widetilde{T_{19}}\right) = \mathcal{A}^{(2)}(G_{19}, T_{19})$ | |
| It results<br>$\left\|\tilde{G}_{16}^{(1)} - \tilde{G}_{i}^{(2)}\right\| \leq \int_0^t (a_{16})^{(2)} \left\|G_{17}^{(1)} - G_{17}^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{(\widehat{M}_{16})^{(2)} s_{(16)}}\, ds_{(16)} +$<br>$\int_0^t \{(a_{16}^{'})^{(2)} \left\|G_{16}^{(1)} - G_{16}^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} +$<br>$(a_{16}^{''})^{(2)}\left(T_{17}^{(1)}, s_{(16)}\right)\left\|G_{16}^{(1)} - G_{16}^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{(\widehat{M}_{16})^{(2)} s_{(16)}} +$<br>$G_{16}^{(2)} \left\|(a_{16}^{''})^{(2)}\left(T_{17}^{(1)}, s_{(16)}\right) - (a_{16}^{''})^{(2)}\left(T_{17}^{(2)}, s_{(16)}\right)\right\| \; e^{-(\widehat{M}_{16})^{(2)} s_{(16)}} e^{(\widehat{M}_{16})^{(2)} s_{(16)}} \} ds_{(16)}$ | 196 |
| Where $s_{(16)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | 197 |
| $\left\|(G_{19})^{(1)} - (G_{19})^{(2)}\right\| e^{-(\widehat{M}_{16})^{(2)} t} \leq$<br>$\frac{1}{(\widehat{M}_{16})^{(2)}}\left((a_{16})^{(2)} + (a_{16}^{'})^{(2)} + (\widehat{A}_{16})^{(2)} + (\widehat{P}_{16})^{(2)}(\widehat{k}_{16})^{(2)}\right) d\left(((G_{19})^{(1)}, (T_{19})^{(1)}; (G_{19})^{(2)}, (T_{19})^{(2)})\right)$ | |
| And analogous inequalities for $G_i$ and $T_i$. Taking into account the hypothesis the result follows | 198 |
| **Remark 6:** The fact that we supposed $(a_{16}^{''})^{(2)}$ and $(b_{16}^{''})^{(2)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{16})^{(2)} e^{(\widehat{M}_{16})^{(2)} t}$ and $(\widehat{Q}_{16})^{(2)} e^{(\widehat{M}_{16})^{(2)} t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}, i = 16,17,18$ depend only on $T_{17}$ and respectively on $(G_{19})$(and not on t) and hypothesis can replaced by a usual Lipschitz condition. | 199 |
| **Remark 7:** There does not exist any t where $G_i(t) = 0$ and $T_i(t) = 0$<br>it results<br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(2)} - (a_i^{''})^{(2)}(T_{17}(s_{(16)}), s_{(16)})\} ds_{(16)}\right]} \geq 0$<br>$T_i(t) \geq T_i^0 e^{(-(b_i^{'})^{(2)} t)} > 0$ for $t > 0$ | 200 |
| **Definition of** $\left((\widehat{M}_{16})^{(2)}\right)_1$, $\left((\widehat{M}_{16})^{(2)}\right)_2$ and $\left((\widehat{M}_{16})^{(2)}\right)_3$ :<br>**Remark 8:** if $G_{16}$ is bounded, the same property have also $G_{17}$ and $G_{18}$ . indeed if<br>$G_{16} < (\widehat{M}_{16})^{(2)}$ it follows $\frac{dG_{17}}{dt} \leq \left((\widehat{M}_{16})^{(2)}\right)_1 - (a_{17}^{'})^{(2)} G_{17}$ and by integrating<br>$G_{17} \leq \left((\widehat{M}_{16})^{(2)}\right)_2 = G_{17}^0 + 2(a_{17})^{(2)}\left((\widehat{M}_{16})^{(2)}\right)_1 / (a_{17}^{'})^{(2)}$ | 201 |

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| In the same way , one can obtain<br>$G_{18} \leq \left((\widehat{M}_{16})^{(2)}\right)_3 = G_{18}^0 + 2(a_{18})^{(2)}\left((\widehat{M}_{16})^{(2)}\right)_2/(a_{18}')^{(2)}$<br>If $G_{17}$ or $G_{18}$ is bounded, the same property follows for $G_{16}$ , $G_{18}$ and $G_{16}$ , $G_{17}$ respectively. | |
| **Remark 9:** If $G_{16}$ is bounded, from below, the same property holds for $G_{17}$ and $G_{18}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{17}$ is bounded from below. | 202 |
| **Remark 10:** If $T_{16}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(2)}((G_{19})(t),t)) = (b_{17}')^{(2)}$ then $T_{17} \to \infty$.<br>**Definition of** $(m)^{(2)}$ and $\varepsilon_2$ :<br>Indeed let $t_2$ be so that for $t > t_2$<br>$(b_{17})^{(2)} - (b_i'')^{(2)}((G_{19})(t),t) < \varepsilon_2, T_{16}(t) > (m)^{(2)}$ | 203 |
| Then $\frac{dT_{17}}{dt} \geq (a_{17})^{(2)}(m)^{(2)} - \varepsilon_2 T_{17}$ which leads to<br>$T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{\varepsilon_2}\right)(1 - e^{-\varepsilon_2 t}) + T_{17}^0 e^{-\varepsilon_2 t}$ If we take t such that $e^{-\varepsilon_2 t} = \frac{1}{2}$ it results | 204 |
| $T_{17} \geq \left(\frac{(a_{17})^{(2)}(m)^{(2)}}{2}\right)$, $t = \log\frac{2}{\varepsilon_2}$ By taking now $\varepsilon_2$ sufficiently small one sees that $T_{17}$ is unbounded.<br>The same property holds for $T_{18}$ if $\lim_{t\to\infty}(b_{18}'')^{(2)}\left((G_{19})(t),t\right) = (b_{18}')^{(2)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 205 |
| It is now sufficient to take $\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}$ , $\frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}} < 1$ and to choose<br>$(\hat{P}_{20})^{(3)}$ and $(\hat{Q}_{20})^{(3)}$ large to have | 207 |
| $\frac{(a_i)^{(3)}}{(\hat{M}_{20})^{(3)}}\left[(\hat{P}_{20})^{(3)} + \left((\hat{P}_{20})^{(3)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{20})^{(3)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{20})^{(3)}$ | 208 |
| $\frac{(b_i)^{(3)}}{(\hat{M}_{20})^{(3)}}\left[\left((\hat{Q}_{20})^{(3)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{20})^{(3)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{20})^{(3)}\right] \leq (\hat{Q}_{20})^{(3)}$ | 209 |
| In order that the operator $\mathcal{A}^{(3)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying Equations into itself | 210 |
| The operator $\mathcal{A}^{(3)}$ is a contraction with respect to the metric<br>$d\left(((G_{23})^{(1)},(T_{23})^{(1)}),((G_{23})^{(2)},(T_{23})^{(2)})\right) =$ | 211 |

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| $\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{20})^{(3)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{20})^{(3)}t}\}$ | |
| Indeed if we denote<br>**Definition of** $\widetilde{G_{23}},\widetilde{T_{23}}:(\widetilde{(G_{23})},\widetilde{(T_{23})})=\mathcal{A}^{(3)}((G_{23}),(T_{23}))$ | 212 |
| It results<br>$\left\|\tilde{G}_{20}^{(1)}-\tilde{G}_i^{(2)}\right\|\le\int_0^t(a_{20})^{(3)}\left\|G_{21}^{(1)}-G_{21}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{(\widehat{M}_{20})^{(3)}s_{(20)}}\,ds_{(20)}+$<br>$\int_0^t\{(a_{20}^{'})^{(3)}\left\|G_{20}^{(1)}-G_{20}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}+$<br>$(a_{20}^{''})^{(3)}\left(T_{21}^{(1)},s_{(20)}\right)\left\|G_{20}^{(1)}-G_{20}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{(\widehat{M}_{20})^{(3)}s_{(20)}}+$<br>$G_{20}^{(2)}\left\|(a_{20}^{''})^{(3)}\left(T_{21}^{(1)},s_{(20)}\right)-(a_{20}^{''})^{(3)}\left(T_{21}^{(2)},s_{(20)}\right)\right\|\ e^{-(\widehat{M}_{20})^{(3)}s_{(20)}}e^{(\widehat{M}_{20})^{(3)}s_{(20)}}\}ds_{(20)}$<br>Where $s_{(20)}$ represents integrand that is integrated over the interval $[0,t]$<br>From the hypotheses it follows | 213 |
| $\left\|G_{23}{}^{(1)}-G_{23}{}^{(2)}\right\|e^{-(\widehat{M}_{20})^{(3)}t}\le$<br>$\frac{1}{(\widehat{M}_{20})^{(3)}}\left((a_{20})^{(3)}+(a_{20}^{'})^{(3)}+(\widehat{A}_{20})^{(3)}+(\widehat{P}_{20})^{(3)}(\widehat{k}_{20})^{(3)}\right)d\left(((G_{23})^{(1)},(T_{23})^{(1)};\ (G_{23})^{(2)},(T_{23})^{(2)})\right)$<br>And analogous inequalities for $G_i\ and\ T_i$. Taking into account the hypothesis the result follows | 214 |
| **Remark 11:** The fact that we supposed $(a_{20}^{''})^{(3)}$ and $(b_{20}^{''})^{(3)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}\ and\ (\widehat{Q}_{20})^{(3)}e^{(\widehat{M}_{20})^{(3)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(3)}$ and $(b_i^{''})^{(3)}, i=20,21,22$ depend only on $T_{21}$ and respectively on $(G_{23})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 215 |
| **Remark 12:** There does not exist any $t$ where $G_i(t)=0\ and\ T_i(t)=0$<br>it results<br>$G_i(t)\ge G_i^0e^{\left[-\int_0^t\{(a_i^{'})^{(3)}-(a_i^{''})^{(3)}(T_{21}(s_{(20)}),s_{(20)})\}ds_{(20)}\right]}\ge 0$<br>$T_i(t)\ge T_i^0e^{(-(b_i^{'})^{(3)}t)}>0$ for t > 0 | 216 |

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| **Definition of** $\left((\widehat{M}_{20})^{(3)}\right)_1, \left((\widehat{M}_{20})^{(3)}\right)_2$ *and* $\left((\widehat{M}_{20})^{(3)}\right)_3$ : | 217 |
| **Remark 13:** if $G_{20}$ is bounded, the same property have also $G_{21}$ *and* $G_{22}$ . indeed if<br>$G_{20} < (\widehat{M}_{20})^{(3)}$ it follows $\frac{dG_{21}}{dt} \leq \left((\widehat{M}_{20})^{(3)}\right)_1 - (a_{21}^{'})^{(3)} G_{21}$ and by integrating<br>$G_{21} \leq \left((\widehat{M}_{20})^{(3)}\right)_2 = G_{21}^0 + 2(a_{21})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_1/(a_{21}^{'})^{(3)}$<br>In the same way , one can obtain<br>$G_{22} \leq \left((\widehat{M}_{20})^{(3)}\right)_3 = G_{22}^0 + 2(a_{22})^{(3)}\left((\widehat{M}_{20})^{(3)}\right)_2/(a_{22}^{'})^{(3)}$<br>If $G_{21}$ *or* $G_{22}$ is bounded, the same property follows for $G_{20}$ , $G_{22}$ and $G_{20}$ , $G_{21}$ respectively. | |
| **Remark 14:** If $G_{20}$ *is* bounded, from below, the same property holds for $G_{21}$ *and* $G_{22}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{21}$ is bounded from below. | 218 |
| **Remark 15:** If $T_{20}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(3)}\left((G_{23})(t),t\right) = (b_{21}^{'})^{(3)}$ then $T_{21} \to \infty$.<br>**Definition of** $(m)^{(3)}$ and $\varepsilon_3$ :<br>Indeed let $t_3$ be so that for $t > t_3$<br>$(b_{21})^{(3)} - (b_i^{''})^{(3)}\left((G_{23})(t),t\right) < \varepsilon_3, T_{20}(t) > (m)^{(3)}$ | 219 |
| Then $\frac{dT_{21}}{dt} \geq (a_{21})^{(3)}(m)^{(3)} - \varepsilon_3 T_{21}$ which leads to<br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{\varepsilon_3}\right)(1 - e^{-\varepsilon_3 t}) + T_{21}^0 e^{-\varepsilon_3 t}$ If we take t such that $e^{-\varepsilon_3 t} = \frac{1}{2}$ it results<br>$T_{21} \geq \left(\frac{(a_{21})^{(3)}(m)^{(3)}}{2}\right)$, $t = log\frac{2}{\varepsilon_3}$ By taking now $\varepsilon_3$ sufficiently small one sees that $T_{21}$ is unbounded.<br>The same property holds for $T_{22}$ if $\lim_{t\to\infty}(b_{22}^{''})^{(3)}\left((G_{23})(t),t\right) = (b_{22}^{'})^{(3)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 220 |
| It is now sufficient to take $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}$ , $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}} < 1$ and to choose<br>$(\widehat{P}_{24})^{(4)}$ and $(\widehat{Q}_{24})^{(4)}$ large to have | 221 |
| $\frac{(a_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[(\widehat{P}_{24})^{(4)} + \left((\widehat{P}_{24})^{(4)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{24})^{(4)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{24})^{(4)}$ | 222 |
| $\frac{(b_i)^{(4)}}{(\widehat{M}_{24})^{(4)}}\left[\left((\widehat{Q}_{24})^{(4)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{24})^{(4)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{24})^{(4)}\right] \leq (\widehat{Q}_{24})^{(4)}$ | 223 |

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| In order that the operator $\mathcal{A}^{(4)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | 224 |
| The operator $\mathcal{A}^{(4)}$ is a contraction with respect to the metric<br>$d\left(((G_{27})^{(1)},(T_{27})^{(1)}),((G_{27})^{(2)},(T_{27})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{24})^{(4)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{24})^{(4)}t}\}$<br>Indeed if we denote<br>**Definition of** $(\widetilde{G_{27}}),(\widetilde{T_{27}})$ : $\left((\widetilde{G_{27}}),(\widetilde{T_{27}})\right)=\mathcal{A}^{(4)}((G_{27}),(T_{27}))$<br>It results<br>$\left\|\tilde{G}_{24}^{(1)}-\tilde{G}_i^{(2)}\right\|\le\int_0^t(a_{24})^{(4)}\left\|G_{25}^{(1)}-G_{25}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}\,ds_{(24)}+$<br>$\int_0^t\{(a_{24}^{'})^{(4)}\left\|G_{24}^{(1)}-G_{24}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}+$<br>$(a_{24}^{''})^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)\left\|G_{24}^{(1)}-G_{24}^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}+$<br>$G_{24}^{(2)}\left\|(a_{24}^{''})^{(4)}\left(T_{25}^{(1)},s_{(24)}\right)-(a_{24}^{''})^{(4)}\left(T_{25}^{(2)},s_{(24)}\right)\right\|\ e^{-(\widehat{M}_{24})^{(4)}s_{(24)}}e^{(\widehat{M}_{24})^{(4)}s_{(24)}}\}ds_{(24)}$<br>Where $s_{(24)}$ represents integrand that is integrated over the interval $[0,\mathrm{t}]$<br>From the hypotheses on Equations it follows | 225 |
| $\left\|(G_{27})^{(1)}-(G_{27})^{(2)}\right\|e^{-(\widehat{M}_{24})^{(4)}t}\le$<br>$\frac{1}{(\widehat{M}_{24})^{(4)}}\left((a_{24})^{(4)}+(a_{24}^{'})^{(4)}+(\widehat{A}_{24})^{(4)}+(\widehat{P}_{24})^{(4)}(\widehat{k}_{24})^{(4)}\right)d\left(((G_{27})^{(1)},(T_{27})^{(1)};\ (G_{27})^{(2)},(T_{27})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 226 |
| **Remark 16:** The fact that we supposed $(a_{24}^{''})^{(4)}$ and $(b_{24}^{''})^{(4)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}$ *and* $(\widehat{Q}_{24})^{(4)}e^{(\widehat{M}_{24})^{(4)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(4)}$ and $(b_i^{''})^{(4)}, i=24{,}25{,}26$ depend only on $\mathrm{T}_{25}$ and respectively on $(G_{27})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 227 |
| **Remark 17:** There does not exist any $t$ where $G_i(t)=0$ *and* $T_i(t)=0$<br>it results | 228 |

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| $G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i')^{(4)} - (a_i'')^{(4)}(T_{25}(s_{(24)}), s_{(24)})\} ds_{(24)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{\left(-(b_i')^{(4)} t\right)} > 0$ for t > 0 | |
| **Definition of** $\left((\widehat{M}_{24})^{(4)}\right)_1, \left((\widehat{M}_{24})^{(4)}\right)_2$ *and* $\left((\widehat{M}_{24})^{(4)}\right)_3$ :<br><br>**Remark 18:** if $G_{24}$ is bounded, the same property have also $G_{25}$ *and* $G_{26}$ . indeed if<br><br>$G_{24} < (\widehat{M}_{24})^{(4)}$ it follows $\frac{dG_{25}}{dt} \leq \left((\widehat{M}_{24})^{(4)}\right)_1 - (a_{25}')^{(4)} G_{25}$ and by integrating<br><br>$G_{25} \leq \left((\widehat{M}_{24})^{(4)}\right)_2 = G_{25}^0 + 2(a_{25})^{(4)} \left((\widehat{M}_{24})^{(4)}\right)_1 / (a_{25}')^{(4)}$<br><br>In the same way , one can obtain<br><br>$G_{26} \leq \left((\widehat{M}_{24})^{(4)}\right)_3 = G_{26}^0 + 2(a_{26})^{(4)} \left((\widehat{M}_{24})^{(4)}\right)_2 / (a_{26}')^{(4)}$<br><br>If $G_{25}$ *or* $G_{26}$ is bounded, the same property follows for $G_{24}$ , $G_{26}$ and $G_{24}$ , $G_{25}$ respectively. | 229 |
| **Remark 19:** If $G_{24}$ *is* bounded, from below, the same property holds for $G_{25}$ *and* $G_{26}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{25}$ is bounded from below. | 230 |
| **Remark 20:** If $T_{24}$ is bounded from below and $\lim_{t\to\infty}((b_i'')^{(4)}((G_{27})(t), t)) = (b_{25}')^{(4)}$ then $T_{25} \to \infty$.<br><br>**Definition of** $(m)^{(4)}$ and $\varepsilon_4$ :<br><br>Indeed let $t_4$ be so that for $t > t_4$<br><br>$(b_{25})^{(4)} - (b_i'')^{(4)}((G_{27})(t), t) < \varepsilon_4, T_{24}(t) > (m)^{(4)}$ | 231 |
| Then $\frac{dT_{25}}{dt} \geq (a_{25})^{(4)}(m)^{(4)} - \varepsilon_4 T_{25}$ which leads to<br><br>$T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{\varepsilon_4}\right)(1 - e^{-\varepsilon_4 t}) + T_{25}^0 e^{-\varepsilon_4 t}$ If we take t such that $e^{-\varepsilon_4 t} = \frac{1}{2}$ it results<br><br>$T_{25} \geq \left(\frac{(a_{25})^{(4)}(m)^{(4)}}{2}\right)$, $t = log\frac{2}{\varepsilon_4}$ By taking now $\varepsilon_4$ sufficiently small one sees that $T_{25}$ is unbounded. The same property holds for $T_{26}$ if $\lim_{t\to\infty}(b_{26}'')^{(4)}\left((G_{27})(t), t\right) = (b_{26}')^{(4)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 42<br><br>Analogous inequalities hold also for $G_{29}$ , $G_{30}, T_{28}, T_{29}, T_{30}$ | 232 |
| It is now sufficient to take $\frac{(a_i)^{(5)}}{(\widehat{M}_{28})^{(5)}}$ , $\frac{(b_i)^{(5)}}{(\widehat{M}_{28})^{(5)}} < 1$ and to choose<br><br>$(\widehat{P}_{28})^{(5)}$ and $(\widehat{Q}_{28})^{(5)}$ large to have | 233 |

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| $\frac{(a_i)^{(5)}}{(\widehat{M}_{28})^{(5)}}\left[(\widehat{P}_{28})^{(5)}+\left((\hat{P}_{28})^{(5)}+G_j^0\right)e^{-\left(\frac{(\hat{P}_{28})^{(5)}+G_j^0}{G_j^0}\right)}\right]\leq(\hat{P}_{28})^{(5)}$ | 234 |
| $\frac{(b_i)^{(5)}}{(\widehat{M}_{28})^{(5)}}\left[\left((\hat{Q}_{28})^{(5)}+T_j^0\right)e^{-\left(\frac{(\hat{Q}_{28})^{(5)}+T_j^0}{T_j^0}\right)}+(\hat{Q}_{28})^{(5)}\right]\leq(\hat{Q}_{28})^{(5)}$ | 235 |
| In order that the operator $\mathcal{A}^{(5)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(5)}$ is a contraction with respect to the metric<br>$d\left(((G_{31})^{(1)},(T_{31})^{(1)}),((G_{31})^{(2)},(T_{31})^{(2)})\right)=$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t)-G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{28})^{(5)}t},\max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t)-T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{28})^{(5)}t}\}$<br>Indeed if we denote<br>**Definition of** $\widetilde{(G_{31})},\widetilde{(T_{31})}$ : $\left(\widetilde{(G_{31})},\widetilde{(T_{31})}\right)=\mathcal{A}^{(5)}((G_{31}),(T_{31}))$<br>It results<br>$\left\|\tilde{G}_{28}^{(1)}-\tilde{G}_i^{(2)}\right\|\leq\int_0^t(a_{28})^{(5)}\left\|G_{29}^{(1)}-G_{29}^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{(\widehat{M}_{28})^{(5)}s_{(28)}}ds_{(28)}+$<br>$\int_0^t\{(a_{28}')^{(5)}\left\|G_{28}^{(1)}-G_{28}^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}+$<br>$(a_{28}'')^{(5)}\left(T_{29}^{(1)},s_{(28)}\right)\left\|G_{28}^{(1)}-G_{28}^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{(\widehat{M}_{28})^{(5)}s_{(28)}}+$<br>$G_{28}^{(2)}\left\|(a_{28}'')^{(5)}\left(T_{29}^{(1)},s_{(28)}\right)-(a_{28}'')^{(5)}\left(T_{29}^{(2)},s_{(28)}\right)\right\|\ e^{-(\widehat{M}_{28})^{(5)}s_{(28)}}e^{(\widehat{M}_{28})^{(5)}s_{(28)}}\}ds_{(28)}$<br>Where $s_{(28)}$ represents integrand that is integrated over the interval $[0,t]$<br>From the hypotheses on it follows | 236 |
| $\left\|(G_{31})^{(1)}-(G_{31})^{(2)}\right\|e^{-(\widehat{M}_{28})^{(5)}t}\leq$<br>$\frac{1}{(\widehat{M}_{28})^{(5)}}\left((a_{28})^{(5)}+(a_{28}')^{(5)}+(\widehat{A}_{28})^{(5)}+(\widehat{P}_{28})^{(5)}(\widehat{k}_{28})^{(5)}\right)d\left(((G_{31})^{(1)},(T_{31})^{(1)};\ (G_{31})^{(2)},(T_{31})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 237 |
| **Remark 21:** The fact that we supposed $(a_{28}'')^{(5)}$ and $(b_{28}'')^{(5)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{28})^{(5)}e^{(\widehat{M}_{28})^{(5)}t}$ *and* $(\widehat{Q}_{28})^{(5)}e^{(\widehat{M}_{28})^{(5)}t}$ respectively of $\mathbb{R}_+$.<br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it | 238 |

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| suffices to consider that $(a_i^{''})^{(5)}$ and $(b_i^{''})^{(5)}, i = 28,29,30$ depend only on $T_{29}$ and respectively on $(G_{31})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 22:** There does not exist any $t$ where $G_i\,(t) = 0\ and\ T_i\,(t) = 0$<br><br>it results<br><br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(5)} - (a_i^{''})^{(5)}(T_{29}(s_{(28)}), s_{(28)})\} ds_{(28)}\right]} \geq 0$<br><br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(5)} t\right)} > 0$ for t > 0 | 239 |
| **Definition of** $\left((\widehat{M}_{28})^{(5)}\right)_1, \left((\widehat{M}_{28})^{(5)}\right)_2\ and\ \left((\widehat{M}_{28})^{(5)}\right)_3$ :<br><br>**Remark 23:** if $G_{28}$ is bounded, the same property have also $G_{29}\ and\ G_{30}$ . indeed if<br><br>$G_{28} < (\widehat{M}_{28})^{(5)}$ it follows $\frac{dG_{29}}{dt} \leq \left((\widehat{M}_{28})^{(5)}\right)_1 - (a_{29}^{'})^{(5)} G_{29}$ and by integrating<br><br>$G_{29} \leq \left((\widehat{M}_{28})^{(5)}\right)_2 = G_{29}^0 + 2(a_{29})^{(5)}\left((\widehat{M}_{28})^{(5)}\right)_1 / (a_{29}^{'})^{(5)}$<br><br>In the same way , one can obtain<br><br>$G_{30} \leq \left((\widehat{M}_{28})^{(5)}\right)_3 = G_{30}^0 + 2(a_{30})^{(5)}\left((\widehat{M}_{28})^{(5)}\right)_2 / (a_{30}^{'})^{(5)}$<br><br>If $G_{29}\ or\ G_{30}$ is bounded, the same property follows for $G_{28}$ , $G_{30}$ and $G_{28}$ , $G_{29}$ respectively. | 240 |
| **Remark 24:** If $G_{28}$ *is* bounded, from below, the same property holds for $G_{29}\ and\ G_{30}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{29}$ is bounded from below. | 241 |
| **Remark 25:** If $T_{28}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(5)}((G_{31})(t), t)) = (b_{29}^{'})^{(5)}$ then $T_{29} \to \infty$.<br><br>**Definition of** $(m)^{(5)}$ and $\varepsilon_5$ :<br><br>Indeed let $t_5$ be so that for $t > t_5$<br><br>$(b_{29})^{(5)} - (b_i^{''})^{(5)}((G_{31})(t), t) < \varepsilon_5, T_{28}\,(t) > (m)^{(5)}$ | 242 |
| Then $\frac{dT_{29}}{dt} \geq (a_{29})^{(5)}(m)^{(5)} - \varepsilon_5 T_{29}$ which leads to<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{\varepsilon_5}\right)(1 - e^{-\varepsilon_5 t}) + T_{29}^0 e^{-\varepsilon_5 t}$ If we take t such that $e^{-\varepsilon_5 t} = \frac{1}{2}$ it results<br><br>$T_{29} \geq \left(\frac{(a_{29})^{(5)}(m)^{(5)}}{2}\right), \quad t = log\frac{2}{\varepsilon_5}$ By taking now $\varepsilon_5$ sufficiently small one sees that $T_{29}$ is unbounded.<br>The same property holds for $T_{30}$ if $\lim_{t\to\infty}(b_{30}^{''})^{(5)}\left((G_{31})(t), t\right) = (b_{30}^{'})^{(5)}$<br><br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations<br><br>Analogous inequalities hold also for $G_{33}$ , $G_{34}, T_{32}, T_{33}, T_{34}$ | 243 |

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| It is now sufficient to take $\frac{(a_i)^{(6)}}{(\widehat{M}_{32})^{(6)}}$ , $\frac{(b_i)^{(6)}}{(\widehat{M}_{32})^{(6)}} < 1$ and to choose<br>$(\widehat{P}_{32})^{(6)}$ and $(\widehat{Q}_{32})^{(6)}$ large to have | 244 |
| $\frac{(a_i)^{(6)}}{(\widehat{M}_{32})^{(6)}}\left[(\widehat{P}_{32})^{(6)} + \left((\widehat{P}_{32})^{(6)} + G_j^0\right)e^{-\left(\frac{(\widehat{P}_{32})^{(6)}+G_j^0}{G_j^0}\right)}\right] \leq (\widehat{P}_{32})^{(6)}$ | 245 |
| $\frac{(b_i)^{(6)}}{(\widehat{M}_{32})^{(6)}}\left[\left((\widehat{Q}_{32})^{(6)} + T_j^0\right)e^{-\left(\frac{(\widehat{Q}_{32})^{(6)}+T_j^0}{T_j^0}\right)} + (\widehat{Q}_{32})^{(6)}\right] \leq (\widehat{Q}_{32})^{(6)}$ | 246 |
| In order that the operator $\mathcal{A}^{(6)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(6)}$ is a contraction with respect to the metric<br>$d\left(((G_{35})^{(1)}, (T_{35})^{(1)}), ((G_{35})^{(2)}, (T_{35})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\widehat{M}_{32})^{(6)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\widehat{M}_{32})^{(6)}t}\}$<br>Indeed if we denote<br>**Definition of** $\widetilde{(G_{35})}, \widetilde{(T_{35})}$ : $\left(\widetilde{(G_{35})}, \widetilde{(T_{35})}\right) = \mathcal{A}^{(6)}((G_{35}), (T_{35}))$<br>It results<br>$\left\|\tilde{G}_{32}^{(1)} - \tilde{G}_i^{(2)}\right\| \leq \int_0^t (a_{32})^{(6)}\left\|G_{33}^{(1)} - G_{33}^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{(\widehat{M}_{32})^{(6)}s_{(32)}}\,ds_{(32)} +$<br>$\int_0^t\{(a'_{32})^{(6)}\left\|G_{32}^{(1)} - G_{32}^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{-(\widehat{M}_{32})^{(6)}s_{(32)}} +$<br>$(a''_{32})^{(6)}\left(T_{33}^{(1)}, s_{(32)}\right)\left\|G_{32}^{(1)} - G_{32}^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{(\widehat{M}_{32})^{(6)}s_{(32)}} +$<br>$G_{32}^{(2)}\left\|(a''_{32})^{(6)}\left(T_{33}^{(1)}, s_{(32)}\right) - (a''_{32})^{(6)}\left(T_{33}^{(2)}, s_{(32)}\right)\right\| e^{-(\widehat{M}_{32})^{(6)}s_{(32)}}e^{(\widehat{M}_{32})^{(6)}s_{(32)}}\}ds_{(32)}$<br>Where $s_{(32)}$ represents integrand that is integrated over the interval $[0, t]$<br>From the hypotheses it follows | 247 |
| $\left\|(G_{35})^{(1)} - (G_{35})^{(2)}\right\|e^{-(\widehat{M}_{32})^{(6)}t} \leq$<br>$\frac{1}{(\widehat{M}_{32})^{(6)}}\left((a_{32})^{(6)} + (a'_{32})^{(6)} + (\widehat{A}_{32})^{(6)} + (\widehat{P}_{32})^{(6)}(\widehat{k}_{32})^{(6)}\right)d\left(((G_{35})^{(1)}, (T_{35})^{(1)}; (G_{35})^{(2)}, (T_{35})^{(2)})\right)$<br>And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis the result follows | 248 |
| **Remark 26:** The fact that we supposed $(a''_{32})^{(6)}$ and $(b''_{32})^{(6)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate | 249 |

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| condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{32})^{(6)}e^{(\widehat{M}_{32})^{(6)}t}$ *and* $(\widehat{Q}_{32})^{(6)}e^{(\widehat{M}_{32})^{(6)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(6)}$ and $(b_i^{''})^{(6)}, i = 32,33,34$ depend only on $T_{33}$ and respectively on $(G_{35})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | |
| **Remark 27:** There does not exist any $t$ where $G_i(t) = 0\ and\ T_i(t) = 0$<br><br>it results<br><br>$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a_i^{'})^{(6)} - (a_i^{''})^{(6)}(T_{33}(s_{(32)}), s_{(32)})\} ds_{(32)}\right]} \geq 0$<br><br>$T_i(t) \geq T_i^0 e^{(-(b_i^{'})^{(6)}t)} > 0$ for t > 0 | 250 |
| **Definition of** $\left((\widehat{M}_{32})^{(6)}\right)_1, \left((\widehat{M}_{32})^{(6)}\right)_2\ and\ \left((\widehat{M}_{32})^{(6)}\right)_3$ :<br><br>**Remark 28:** if $G_{32}$ is bounded, the same property have also $G_{33}\ and\ G_{34}$ . indeed if<br><br>$G_{32} < (\widehat{M}_{32})^{(6)}$ it follows $\frac{dG_{33}}{dt} \leq \left((\widehat{M}_{32})^{(6)}\right)_1 - (a_{33}^{'})^{(6)}G_{33}$ and by integrating<br><br>$G_{33} \leq \left((\widehat{M}_{32})^{(6)}\right)_2 = G_{33}^0 + 2(a_{33})^{(6)}\left((\widehat{M}_{32})^{(6)}\right)_1/(a_{33}^{'})^{(6)}$<br><br>In the same way , one can obtain<br><br>$G_{34} \leq \left((\widehat{M}_{32})^{(6)}\right)_3 = G_{34}^0 + 2(a_{34})^{(6)}\left((\widehat{M}_{32})^{(6)}\right)_2/(a_{34}^{'})^{(6)}$<br><br>If $G_{33}\ or\ G_{34}$ is bounded, the same property follows for $G_{32}$ , $G_{34}$ and $G_{32}$ , $G_{33}$ respectively. | 251 |
| **Remark 29:** If $G_{32}$ *is* bounded, from below, the same property holds for $G_{33}\ and\ G_{34}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{33}$ is bounded from below. | 252 |
| **Remark 30:** If $T_{32}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(6)}((G_{35})(t), t)) = (b_{33}^{'})^{(6)}$ then $T_{33} \to \infty$.<br><br>**Definition of** $(m)^{(6)}$ and $\varepsilon_6$ :<br><br>Indeed let $t_6$ be so that for $t > t_6$<br><br>$(b_{33})^{(6)} - (b_i^{''})^{(6)}((G_{35})(t), t) < \varepsilon_6, T_{32}(t) > (m)^{(6)}$ | 253 |
| Then $\frac{dT_{33}}{dt} \geq (a_{33})^{(6)}(m)^{(6)} - \varepsilon_6 T_{33}$ which leads to<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{\varepsilon_6}\right)(1 - e^{-\varepsilon_6 t}) + T_{33}^0 e^{-\varepsilon_6 t}$ If we take t such that $e^{-\varepsilon_6 t} = \frac{1}{2}$ it results<br><br>$T_{33} \geq \left(\frac{(a_{33})^{(6)}(m)^{(6)}}{2}\right)$, $t = log\frac{2}{\varepsilon_6}$ By taking now $\varepsilon_6$ sufficiently small one sees that $T_{33}$ is unbounded.<br>The same property holds for $T_{34}$ if $\lim_{t\to\infty}(b_{34}^{''})^{(6)}\left((G_{35})(t), t(t), t\right) = (b_{34}^{'})^{(6)}$ | 254 |

We now state a more precise theorem about the behaviors at infinity of the solutions of equations

Analogous inequalities hold also for $G_{37}, G_{38}, T_{36}, T_{37}, T_{38}$ 255

It is now sufficient to take $\frac{(a_i)^{(7)}}{(\hat{M}_{36})^{(7)}}, \frac{(b_i)^{(7)}}{(\hat{M}_{36})^{(7)}} < 1$ and to choose
$(\hat{P}_{36})^{(7)}$ and $(\hat{Q}_{36})^{(7)}$ large to have

$$\frac{(a_i)^{(7)}}{(\hat{M}_{36})^{(7)}}\left[(\hat{P}_{36})^{(7)} + \left((\hat{P}_{36})^{(7)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{36})^{(7)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{36})^{(7)} \quad 256$$

$$\frac{(b_i)^{(7)}}{(\hat{M}_{36})^{(7)}}\left[\left((\hat{Q}_{36})^{(7)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{36})^{(7)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{36})^{(7)}\right] \leq (\hat{Q}_{36})^{(7)} \quad 257$$

In order that the operator $\mathcal{A}^{(7)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself

The operator $\mathcal{A}^{(7)}$ is a contraction with respect to the metric 258

$$d\left(((G_{39})^{(1)}, (T_{39})^{(1)}), ((G_{39})^{(2)}, (T_{39})^{(2)})\right) =$$
$$\sup_i\{\max_{t\in\mathbb{R}_+}|G_i^{(1)}(t) - G_i^{(2)}(t)|e^{-(\hat{M}_{36})^{(7)}t}, \max_{t\in\mathbb{R}_+}|T_i^{(1)}(t) - T_i^{(2)}(t)|e^{-(\hat{M}_{36})^{(7)}t}\}$$

Indeed if we denote

**Definition of** $(\widetilde{G_{39}}), (\widetilde{T_{39}}) : \left((\widetilde{G_{39}}), (\widetilde{T_{39}})\right) = \mathcal{A}^{(7)}((G_{39}), (T_{39}))$
It results

$$|\tilde{G}_{36}^{(1)} - \tilde{G}_i^{(2)}| \leq \int_0^t (a_{36})^{(7)}|G_{37}^{(1)} - G_{37}^{(2)}|e^{-(\hat{M}_{36})^{(7)}s_{(36)}}e^{(\hat{M}_{36})^{(7)}s_{(36)}}\,ds_{(36)} +$$
$$\int_0^t\{(a'_{36})^{(7)}|G_{36}^{(1)} - G_{36}^{(2)}|e^{-(\hat{M}_{36})^{(7)}s_{(36)}}e^{-(\hat{M}_{36})^{(7)}s_{(36)}} +$$
$$(a''_{36})^{(7)}(T_{37}^{(1)}, s_{(36)})|G_{36}^{(1)} - G_{36}^{(2)}|e^{-(\hat{M}_{36})^{(7)}s_{(36)}}e^{(\hat{M}_{36})^{(7)}s_{(36)}} +$$
$$G_{36}^{(2)}|(a''_{36})^{(7)}(T_{37}^{(1)}, s_{(36)}) - (a''_{36})^{(7)}(T_{37}^{(2)}, s_{(36)})|\ e^{-(\hat{M}_{36})^{(7)}s_{(36)}}e^{(\hat{M}_{36})^{(7)}s_{(36)}}\}ds_{(36)}$$

Where $s_{(36)}$ represents integrand that is integrated over the interval $[0, t]$

From the hypotheses on it follows

$$|(G_{39})^{(1)} - (G_{39})^{(2)}|e^{-(\hat{M}_{36})^{(7)}t} \leq$$
$$\frac{1}{(\hat{M}_{36})^{(7)}}\left((a_{36})^{(7)} + (a'_{36})^{(7)} + (\hat{A}_{36})^{(7)} + (\hat{P}_{36})^{(7)}(\hat{k}_{36})^{(7)}\right)d\left(((G_{39})^{(1)}, (T_{39})^{(1)}; (G_{39})^{(2)}, (T_{39})^{(2)})\right) \quad 259$$

And analogous inequalities for $G_i$ and $T_i$. Taking into account the hypothesis the result follows

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| **Remark 31:** The fact that we supposed $(a_{36}^{''})^{(7)}$ and $(b_{36}^{''})^{(7)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{36})^{(7)}e^{(\widehat{M}_{36})^{(7)}t}$ $and$ $(\widehat{Q}_{36})^{(7)}e^{(\widehat{M}_{36})^{(7)}t}$ respectively of $\mathbb{R}_+$.<br><br>If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a_i^{''})^{(7)}$ and $(b_i^{''})^{(7)}, i = 36,37,38$ depend only on $T_{37}$ and respectively on $(G_{39})(and\ not\ on\ \ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 260 |
| **Remark 32:** There does not exist any $t$ where $G_i\,(t) = 0\ and\ T_i\,(t) = 0$<br><br>it results<br><br>$G_i\,(t) \geq G_i^0 e^{\left[-\int_0^t\{(a_i^{'})^{(7)}-(a_i^{''})^{(7)}(T_{37}(s_{(36)}),s_{(36)})\}ds_{(36)}\right]} \geq 0$<br><br>$T_i\,(t) \geq T_i^0 e^{\left(-(b_i^{'})^{(7)}t\right)} > 0$ for t > 0 | 261 |
| **Definition of** $\left((\widehat{M}_{36})^{(7)}\right)_1, \left((\widehat{M}_{36})^{(7)}\right)_2\ and\ \left((\widehat{M}_{36})^{(7)}\right)_3$ :<br><br>**Remark 33:** if $G_{36}$ is bounded, the same property have also $G_{37}\ and\ G_{38}$ . indeed if<br><br>$G_{36} < (\widehat{M}_{36})^{(7)}$ it follows $\frac{dG_{37}}{dt} \leq \left((\widehat{M}_{36})^{(7)}\right)_1 - (a_{37}^{'})^{(7)}G_{37}$ and by integrating<br><br>$G_{37} \leq \left((\widehat{M}_{36})^{(7)}\right)_2 = G_{37}^0 + 2(a_{37})^{(7)}\left((\widehat{M}_{36})^{(7)}\right)_1/(a_{37}^{'})^{(7)}$<br><br>In the same way , one can obtain<br><br>$G_{38} \leq \left((\widehat{M}_{36})^{(7)}\right)_3 = G_{38}^0 + 2(a_{38})^{(7)}\left((\widehat{M}_{36})^{(7)}\right)_2/(a_{38}^{'})^{(7)}$<br><br>If $G_{37}\ or\ G_{38}$ is bounded, the same property follows for $G_{36}$ , $G_{38}$ and $G_{36}$ , $G_{37}$ respectively. | 262 |
| **Remark 34:** If $G_{36}$ $is$ bounded, from below, the same property holds for $G_{37}\ and\ G_{38}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{37}$ is bounded from below. | 263 |
| **Remark 35:** If $T_{36}$ is bounded from below and $\lim_{t\to\infty}((b_i^{''})^{(7)}\,((G_{39})(t),t)) = (b_{37}^{'})^{(7)}$ then $T_{37} \to \infty$.<br><br>**Definition of** $(m)^{(7)}$ and $\varepsilon_7$ :<br><br>Indeed let $t_7$ be so that for $t > t_7$<br><br>$(b_{37})^{(7)} - (b_i^{''})^{(7)}((G_{39})(t),t) < \varepsilon_7, T_{36}\,(t) > (m)^{(7)}$ | 264 |

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| Then $\frac{dT_{37}}{dt} \geq (a_{37})^{(7)}(m)^{(7)} - \varepsilon_7 T_{37}$ which leads to<br><br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{\varepsilon_7}\right)(1 - e^{-\varepsilon_7 t}) + T_{37}^0 e^{-\varepsilon_7 t}$ If we take t such that $e^{-\varepsilon_7 t} = \frac{1}{2}$ it results<br><br>$T_{37} \geq \left(\frac{(a_{37})^{(7)}(m)^{(7)}}{2}\right), \quad t = log\frac{2}{\varepsilon_7}$ By taking now $\varepsilon_7$ sufficiently small one sees that $T_{37}$ is unbounded.<br>The same property holds for $T_{38}$ if $\lim_{t\to\infty}(b_{38}^{''})^{(7)}\left((G_{39})(t), t\right) = (b_{38}^{'})^{(7)}$<br>We now state a more precise theorem about the behaviors at infinity of the solutions of equations | 265 |
| It is now sufficient to take $\frac{(a_i)^{(8)}}{(\bar{M}_{40})^{(8)}}$ , $\frac{(b_i)^{(8)}}{(\bar{M}_{40})^{(8)}} < 1$ and to choose<br>$(\hat{P}_{40})^{(8)}$ and $(\hat{Q}_{40})^{(8)}$ large to have | 266 |
| $\frac{(a_i)^{(8)}}{(\bar{M}_{40})^{(8)}}\left[(\hat{P}_{40})^{(8)} + \left((\hat{P}_{40})^{(8)} + G_j^0\right)e^{-\left(\frac{(\hat{P}_{40})^{(8)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{40})^{(8)}$ | 267 |
| $\frac{(b_i)^{(8)}}{(\bar{M}_{40})^{(8)}}\left[\left((\hat{Q}_{40})^{(8)} + T_j^0\right)e^{-\left(\frac{(\hat{Q}_{40})^{(8)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{40})^{(8)}\right] \leq (\hat{Q}_{40})^{(8)}$ | 268 |
| In order that the operator $\mathcal{A}^{(8)}$ transforms the space of sextuples of functions $G_i, T_i$ satisfying Equations into itself | |
| The operator $\mathcal{A}^{(8)}$ is a contraction with respect to the metric | |
| $d\left(((G_{43})^{(1)}, (T_{43})^{(1)}), ((G_{43})^{(2)}, (T_{43})^{(2)})\right) =$<br>$\sup_i\{\max_{t\in\mathbb{R}_+}\left\|G_i^{(1)}(t) - G_i^{(2)}(t)\right\|e^{-(\bar{M}_{40})^{(8)}t}, \max_{t\in\mathbb{R}_+}\left\|T_i^{(1)}(t) - T_i^{(2)}(t)\right\|e^{-(\bar{M}_{40})^{(8)}t}\}$ | 269 |
| Indeed if we denote<br><br>**Definition of** $\widetilde{(G_{43})}, \widetilde{(T_{43})}$ : $\left(\widetilde{(G_{43})}, \widetilde{(T_{43})}\right) = \mathcal{A}^{(8)}((G_{43}), (T_{43}))$ | 270 |
| It results<br><br>$\left\|\tilde{G}_{40}^{(1)} - \tilde{G}_i^{(2)}\right\| \leq \int_0^t (a_{40})^{(8)}\left\|G_{41}^{(1)} - G_{41}^{(2)}\right\|e^{-(\bar{M}_{40})^{(8)}s_{(40)}}e^{(\bar{M}_{40})^{(8)}s_{(40)}}\,ds_{(40)} +$<br>$\int_0^t\{(a_{40}^{'})^{(8)}\left\|G_{40}^{(1)} - G_{40}^{(2)}\right\|e^{-(\bar{M}_{40})^{(8)}s_{(40)}}e^{-(\bar{M}_{40})^{(8)}s_{(40)}} +$<br>$(a_{40}^{''})^{(8)}\left(T_{41}^{(1)}, s_{(40)}\right)\left\|G_{40}^{(1)} - G_{40}^{(2)}\right\|e^{-(\bar{M}_{40})^{(8)}s_{(40)}}e^{(\bar{M}_{40})^{(8)}s_{(40)}} +$<br>$G_{40}^{(2)}\left\|(a_{40}^{''})^{(8)}\left(T_{41}^{(1)}, s_{(40)}\right) - (a_{40}^{''})^{(8)}\left(T_{41}^{(2)}, s_{(40)}\right)\right\| \; e^{-(\bar{M}_{40})^{(8)}s_{(40)}}e^{(\bar{M}_{40})^{(8)}s_{(40)}}\}ds_{(40)}$ | 271 |
| Where $s_{(40)}$ represents integrand that is integrated over the interval [0, t] | 272 |

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| From the hypotheses it follows | |
| $\left\|(G_{43})^{(1)} - (G_{43})^{(2)}\right\| e^{-(\widehat{M}_{40})^{(8)}t} \leq$ $\frac{1}{(\widehat{M}_{40})^{(8)}}\left((a_{40})^{(8)} + (a'_{40})^{(8)} + (\widehat{A}_{40})^{(8)} + (\widehat{P}_{40})^{(8)}(\widehat{k}_{40})^{(8)}\right) d\left(((G_{43})^{(1)}, (T_{43})^{(1)};\ (G_{43})^{(2)}, (T_{43})^{(2)})\right)$ And analogous inequalities for $G_i$ $and$ $T_i$. Taking into account the hypothesis the result follows | 273 |
| **Remark 36:** The fact that we supposed $(a''_{40})^{(8)}$ and $(b''_{40})^{(8)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{40})^{(8)} e^{(\widehat{M}_{40})^{(8)}t}$ $and$ $(\widehat{Q}_{40})^{(8)} e^{(\widehat{M}_{40})^{(8)}t}$ respectively of $\mathbb{R}_+$. If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a''_i)^{(8)}$ and $(b''_i)^{(8)}, i = 40,41,42$ depend only on $T_{41}$ and respectively on $(G_{43})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition. | 274 |
| **Remark 37** There does not exist any $t$ where $G_i(t) = 0$ $and$ $T_i(t) = 0$ it results $G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a'_i)^{(8)} - (a''_i)^{(8)}(T_{41}(s_{(40)}), s_{(40)})\} ds_{(40)}\right]} \geq 0$ $T_i(t) \geq T_i^0 e^{\left(-(b'_i)^{(8)}t\right)} > 0$ for t > 0 | 275 |
| **Definition of** $\left((\widehat{M}_{40})^{(8)}\right)_1, \left((\widehat{M}_{40})^{(8)}\right)_2$ $and$ $\left((\widehat{M}_{40})^{(8)}\right)_3$ : **Remark 38:** if $G_{40}$ is bounded, the same property have also $G_{41}$ $and$ $G_{42}$ . indeed if $G_{40} < (\widehat{M}_{40})^{(8)}$ it follows $\frac{dG_{41}}{dt} \leq \left((\widehat{M}_{40})^{(8)}\right)_1 - (a'_{41})^{(8)} G_{41}$ and by integrating $G_{41} \leq \left((\widehat{M}_{40})^{(8)}\right)_2 = G_{41}^0 + 2(a_{41})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_1 / (a'_{41})^{(8)}$ In the same way , one can obtain $G_{42} \leq \left((\widehat{M}_{40})^{(8)}\right)_3 = G_{42}^0 + 2(a_{42})^{(8)}\left((\widehat{M}_{40})^{(8)}\right)_2 / (a'_{42})^{(8)}$ If $G_{41}$ $or$ $G_{42}$ is bounded, the same property follows for $G_{40}$ , $G_{42}$ and $G_{40}$ , $G_{41}$ respectively. | 276 |
| **Remark 39:** If $G_{40}$ $is$ bounded, from below, the same property holds for $G_{41}$ $and$ $G_{42}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{41}$ is bounded from below. | 277 |
| **Remark 40:** If $T_{40}$ is bounded from below and $\lim_{t\to\infty}((b''_i)^{(8)}((G_{43})(t), t)) = (b'_{41})^{(8)}$ then $T_{41} \to \infty$. | 278 |

**Definition of** $(m)^{(8)}$ and $\varepsilon_8$ :

Indeed let $t_8$ be so that for $t > t_8$

$(b_{41})^{(8)} - (b_i^{''})^{(8)}\big((G_{43})(t), t\big) < \varepsilon_8, T_{40}(t) > (m)^{(8)}$

Then $\frac{dT_{41}}{dt} \geq (a_{41})^{(8)}(m)^{(8)} - \varepsilon_8 T_{41}$ which leads to

$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{\varepsilon_8}\right)(1 - e^{-\varepsilon_8 t}) + T_{41}^0 e^{-\varepsilon_8 t}$ If we take t such that $e^{-\varepsilon_8 t} = \frac{1}{2}$ it results

$T_{41} \geq \left(\frac{(a_{41})^{(8)}(m)^{(8)}}{2}\right), \quad t = log\frac{2}{\varepsilon_8}$ By taking now $\varepsilon_8$ sufficiently small one sees that $T_{41}$ is unbounded.
The same property holds for $T_{42}$ if $\lim_{t\to\infty}(b_{42}^{''})^{(8)}\big((G_{43})(t), t(t), t\big) = (b_{42}^{'})^{(8)}$

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It is now sufficient to take $\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}, \frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}} < 1$ and to choose $(\widehat{P}_{44})^{(9)}$ and $(\widehat{Q}_{44})^{(9)}$ large to have

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$$\frac{(a_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[(\widehat{P}_{44})^{(9)} + \big((\hat{P}_{44})^{(9)} + G_j^0\big)e^{-\left(\frac{(\hat{P}_{44})^{(9)}+G_j^0}{G_j^0}\right)}\right] \leq (\hat{P}_{44})^{(9)}$$

$$\frac{(b_i)^{(9)}}{(\widehat{M}_{44})^{(9)}}\left[\big((\hat{Q}_{44})^{(9)} + T_j^0\big)e^{-\left(\frac{(\hat{Q}_{44})^{(9)}+T_j^0}{T_j^0}\right)} + (\hat{Q}_{44})^{(9)}\right] \leq (\hat{Q}_{44})^{(9)}$$

In order that the operator $\mathcal{A}^{(9)}$ transforms the space of sextuples of functions $G_i$ , $T_i$ satisfying 39,35,36 into itself

The operator $\mathcal{A}^{(9)}$ is a contraction with respect to the metric

$$d\Big(\big((G_{47})^{(1)}, (T_{47})^{(1)}\big), \big((G_{47})^{(2)}, (T_{47})^{(2)}\big)\Big) =$$
$$\sup_i\{\max_{t\in\mathbb{R}_+}\big|G_i^{(1)}(t) - G_i^{(2)}(t)\big|e^{-(\widehat{M}_{44})^{(9)}t}, \max_{t\in\mathbb{R}_+}\big|T_i^{(1)}(t) - T_i^{(2)}(t)\big|e^{-(\widehat{M}_{44})^{(9)}t}\}$$

Indeed if we denote

**Definition of** $\widetilde{(G_{47})}, \widetilde{(T_{47})} : \big(\widetilde{(G_{47})}, \widetilde{(T_{47})}\big) = \mathcal{A}^{(9)}\big((G_{47}), (T_{47})\big)$
It results

$$\big|\tilde{G}_{44}^{(1)} - \tilde{G}_i^{(2)}\big| \leq \int_0^t (a_{44})^{(9)}\big|G_{45}^{(1)} - G_{45}^{(2)}\big|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\,ds_{(44)} +$$
$$\int_0^t\{(a_{44}^{'})^{(9)}\big|G_{44}^{(1)} - G_{44}^{(2)}\big|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{-(\widehat{M}_{44})^{(9)}s_{(44)}} +$$
$$(a_{44}^{''})^{(9)}\big(T_{45}^{(1)}, s_{(44)}\big)\big|G_{44}^{(1)} - G_{44}^{(2)}\big|e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}} +$$
$$G_{44}^{(2)}\big|(a_{44}^{''})^{(9)}\big(T_{45}^{(1)}, s_{(44)}\big) - (a_{44}^{''})^{(9)}\big(T_{45}^{(2)}, s_{(44)}\big)\big|\ e^{-(\widehat{M}_{44})^{(9)}s_{(44)}}e^{(\widehat{M}_{44})^{(9)}s_{(44)}}\}ds_{(44)}$$

Where $s_{(44)}$ represents integrand that is integrated over the interval $[0, t]$

From the hypotheses on 45,46,47,28 and 29 it follows

$$\left|(G_{47})^{(1)} - G^{(2)}\right| e^{-(\widehat{M}_{44})^{(9)}t} \leq \frac{1}{(\widehat{M}_{44})^{(9)}}\left((a_{44})^{(9)} + (a'_{44})^{(9)} + (\widehat{A}_{44})^{(9)} + (\widehat{P}_{44})^{(9)}(\widehat{k}_{44})^{(9)}\right) d\left(((G_{47})^{(1)}, (T_{47})^{(1)};\ (G_{47})^{(2)}, (T_{47})^{(2)})\right)$$

And analogous inequalities for $G_i$ *and* $T_i$. Taking into account the hypothesis (39,35,36) the result follows

**Remark 41:** The fact that we supposed $(a''_{44})^{(9)}$ and $(b''_{44})^{(9)}$ depending also on t can be considered as not conformal with the reality, however we have put this hypothesis ,in order that we can postulate condition necessary to prove the uniqueness of the solution bounded by $(\widehat{P}_{44})^{(9)} e^{(\widehat{M}_{44})^{(9)}t}$ *and* $(\widehat{Q}_{44})^{(9)} e^{(\widehat{M}_{44})^{(9)}t}$ respectively of $\mathbb{R}_+$.

If instead of proving the existence of the solution on $\mathbb{R}_+$, we have to prove it only on a compact then it suffices to consider that $(a''_i)^{(9)}$ and $(b''_i)^{(9)}, i = 44,45,46$ depend only on $T_{45}$ and respectively on $(G_{47})(and\ not\ on\ t)$ and hypothesis can replaced by a usual Lipschitz condition.

**Remark 42:** There does not exist any $t$ where $G_i(t) = 0\ and\ T_i(t) = 0$

From 99 to 44 it results

$$G_i(t) \geq G_i^0 e^{\left[-\int_0^t \{(a'_i)^{(9)} - (a''_i)^{(9)}(T_{45}(s_{(44)}), s_{(44)})\} ds_{(44)}\right]} \geq 0$$

$$T_i(t) \geq T_i^0 e^{(-(b'_i)^{(9)}t)} > 0 \quad \text{for t} > 0$$

**Definition of** $\left((\widehat{M}_{44})^{(9)}\right)_1, \left((\widehat{M}_{44})^{(9)}\right)_2$ *and* $\left((\widehat{M}_{44})^{(9)}\right)_3$ :

**Remark 43:** if $G_{44}$ is bounded, the same property have also $G_{45}\ and\ G_{46}$ . indeed if

$G_{44} < (\widehat{M}_{44})^{(9)}$ it follows $\frac{dG_{45}}{dt} \leq \left((\widehat{M}_{44})^{(9)}\right)_1 - (a'_{45})^{(9)} G_{45}$ and by integrating

$$G_{45} \leq \left((\widehat{M}_{44})^{(9)}\right)_2 = G_{45}^0 + 2(a_{45})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_1 / (a'_{45})^{(9)}$$

In the same way , one can obtain

$$G_{46} \leq \left((\widehat{M}_{44})^{(9)}\right)_3 = G_{46}^0 + 2(a_{46})^{(9)}\left((\widehat{M}_{44})^{(9)}\right)_2 / (a'_{46})^{(9)}$$

If $G_{45}\ or\ G_{46}$ is bounded, the same property follows for $G_{44}$ , $G_{46}$ and $G_{44}$ , $G_{45}$ respectively.

**Remark 44:** If $G_{44}$ *is* bounded, from below, the same property holds for $G_{45}\ and\ G_{46}$ . The proof is analogous with the preceding one. An analogous property is true if $G_{45}$ is bounded from below.

**Remark 45:** If $T_{44}$ is bounded from below and $\lim_{t\to\infty}((b''_i)^{(9)}\left((G_{47})(t), t\right)) = (b'_{45})^{(9)}$ then $T_{45} \to \infty$.

**Definition of** $(m)^{(9)}$ and $\varepsilon_9$ :

Indeed let $t_9$ be so that for $t > t_9$

$$(b_{45})^{(9)} - (b_i^{''})^{(9)}\big((G_{47})(t),t\big) < \varepsilon_9, T_{44}\,(t) > (m)^{(9)}$$

Then $\frac{dT_{45}}{dt} \geq (a_{45})^{(9)}(m)^{(9)} - \varepsilon_9 T_{45}$ which leads to

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{\varepsilon_9}\right)(1 - e^{-\varepsilon_9 t}) + T_{45}^0 e^{-\varepsilon_9 t}$ If we take t such that $e^{-\varepsilon_9 t} = \frac{1}{2}$ it results

$T_{45} \geq \left(\frac{(a_{45})^{(9)}(m)^{(9)}}{2}\right)$, $t = log\frac{2}{\varepsilon_9}$ By taking now $\varepsilon_9$ sufficiently small one sees that $T_{45}$ is unbounded.
The same property holds for $T_{46}$ if $\lim_{t\to\infty}(b_{46}^{''})^{(9)}\left((G_{47})(t),t\right) = (b_{46}^{'})^{(9)}$
We now state a more precise theorem about the behaviors at infinity of the solutions of equations 37 to 92

**Behavior of the solutions of equation** 280

**Theorem** If we denote and define

**Definition of** $(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ :

$(\sigma_1)^{(1)}, (\sigma_2)^{(1)}, (\tau_1)^{(1)}, (\tau_2)^{(1)}$ four constants satisfying
$-(\sigma_2)^{(1)} \leq -(a_{13}^{'})^{(1)} + (a_{14}^{'})^{(1)} - (a_{13}^{''})^{(1)}(T_{14}, t) + (a_{14}^{''})^{(1)}(T_{14}, t) \leq -(\sigma_1)^{(1)}$

$-(\tau_2)^{(1)} \leq -(b_{13}^{'})^{(1)} + (b_{14}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G,t) - (b_{14}^{''})^{(1)}(G,t) \leq -(\tau_1)^{(1)}$

**Definition of** $(\nu_1)^{(1)}, (\nu_2)^{(1)}, (u_1)^{(1)}, (u_2)^{(1)}, \nu^{(1)}, u^{(1)}$ : 281

By $(\nu_1)^{(1)} > 0$, $(\nu_2)^{(1)} < 0$ and respectively $(u_1)^{(1)} > 0$, $(u_2)^{(1)} < 0$ the roots of the equations
$(a_{14})^{(1)}\left(\nu^{(1)}\right)^2 + (\sigma_1)^{(1)}\nu^{(1)} - (a_{13})^{(1)} = 0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2 + (\tau_1)^{(1)}u^{(1)} - (b_{13})^{(1)} = 0$

**Definition of** $(\bar{\nu}_1)^{(1)},, (\bar{\nu}_2)^{(1)}, (\bar{u}_1)^{(1)}, (\bar{u}_2)^{(1)}$ : 282

By $(\bar{\nu}_1)^{(1)} > 0$, $(\bar{\nu}_2)^{(1)} < 0$ and respectively $(\bar{u}_1)^{(1)} > 0$, $(\bar{u}_2)^{(1)} < 0$ the roots of the equations
$(a_{14})^{(1)}\left(\nu^{(1)}\right)^2 + (\sigma_2)^{(1)}\nu^{(1)} - (a_{13})^{(1)} = 0$ and $(b_{14})^{(1)}\left(u^{(1)}\right)^2 + (\tau_2)^{(1)}u^{(1)} - (b_{13})^{(1)} = 0$

**Definition of** $(m_1)^{(1)}, (m_2)^{(1)}, (\mu_1)^{(1)}, (\mu_2)^{(1)}, (\nu_0)^{(1)}$ :- 283

If we define $(m_1)^{(1)}, (m_2)^{(1)}, (\mu_1)^{(1)}, (\mu_2)^{(1)}$ by
$(m_2)^{(1)} = (\nu_0)^{(1)}, (m_1)^{(1)} = (\nu_1)^{(1)}$, $if\ (\nu_0)^{(1)} < (\nu_1)^{(1)}$

$(m_2)^{(1)} = (\nu_1)^{(1)}, (m_1)^{(1)} = (\bar{\nu}_1)^{(1)}, if\ (\nu_1)^{(1)} < (\nu_0)^{(1)} < (\bar{\nu}_1)^{(1)}$,

and $\boxed{(\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}}$

$(m_2)^{(1)} = (\nu_1)^{(1)}, (m_1)^{(1)} = (\nu_0)^{(1)}$, $if\ (\bar{\nu}_1)^{(1)} < (\nu_0)^{(1)}$

and analogously 284

$(\mu_2)^{(1)} = (u_0)^{(1)}, (\mu_1)^{(1)} = (u_1)^{(1)}$, $if\ (u_0)^{(1)} < (u_1)^{(1)}$

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| $(\mu_2)^{(1)} = (u_1)^{(1)}, (\mu_1)^{(1)} = (\bar{u}_1)^{(1)}, if\ (u_1)^{(1)} < (u_0)^{(1)} < (\bar{u}_1)^{(1)},$ and $\boxed{(u_0)^{(1)} = \frac{T_{13}^0}{T_{14}^0}}$ $(\mu_2)^{(1)} = (u_1)^{(1)}, (\mu_1)^{(1)} = (u_0)^{(1)}, if\ (\bar{u}_1)^{(1)} < (u_0)^{(1)}$ where $(u_1)^{(1)}, (\bar{u}_1)^{(1)}$ are defined | |
| Then the solution of global equations satisfies the inequalities $G_{13}^0 e^{((S_1)^{(1)} - (p_{13})^{(1)})t} \le G_{13}(t) \le G_{13}^0 e^{(S_1)^{(1)} t}$ where $(p_i)^{(1)}$ is defined by equation $\frac{1}{(m_1)^{(1)}} G_{13}^0 e^{((S_1)^{(1)} - (p_{13})^{(1)})t} \le G_{14}(t) \le \frac{1}{(m_2)^{(1)}} G_{13}^0 e^{(S_1)^{(1)} t}$ | 285 |
| $\left(\frac{(a_{15})^{(1)} G_{13}^0}{(m_1)^{(1)}((S_1)^{(1)} - (p_{13})^{(1)} - (S_2)^{(1)})}\left[e^{((S_1)^{(1)} - (p_{13})^{(1)})t} - e^{-(S_2)^{(1)} t}\right] + G_{15}^0 e^{-(S_2)^{(1)} t} \le G_{15}(t) \le \frac{(a_{15})^{(1)} G_{13}^0}{(m_2)^{(1)}((S_1)^{(1)} - (a'_{15})^{(1)})}[e^{(S_1)^{(1)} t} - e^{-(a'_{15})^{(1)} t}] + G_{15}^0 e^{-(a'_{15})^{(1)} t}\right)$ | 286 |
| $\boxed{T_{13}^0 e^{(R_1)^{(1)} t} \le T_{13}(t) \le T_{13}^0 e^{((R_1)^{(1)} + (r_{13})^{(1)})t}}$ | 287 |
| $\frac{1}{(\mu_1)^{(1)}} T_{13}^0 e^{(R_1)^{(1)} t} \le T_{13}(t) \le \frac{1}{(\mu_2)^{(1)}} T_{13}^0 e^{((R_1)^{(1)} + (r_{13})^{(1)})t}$ | 288 |
| $\frac{(b_{15})^{(1)} T_{13}^0}{(\mu_1)^{(1)}((R_1)^{(1)} - (b'_{15})^{(1)})}\left[e^{(R_1)^{(1)} t} - e^{-(b'_{15})^{(1)} t}\right] + T_{15}^0 e^{-(b'_{15})^{(1)} t} \le T_{15}(t) \le \frac{(a_{15})^{(1)} T_{13}^0}{(\mu_2)^{(1)}((R_1)^{(1)} + (r_{13})^{(1)} + (R_2)^{(1)})}\left[e^{((R_1)^{(1)} + (r_{13})^{(1)})t} - e^{-(R_2)^{(1)} t}\right] + T_{15}^0 e^{-(R_2)^{(1)} t}$ | 289 |
| **Definition of** $(S_1)^{(1)}, (S_2)^{(1)}, (R_1)^{(1)}, (R_2)^{(1)}$:- Where $(S_1)^{(1)} = (a_{13})^{(1)}(m_2)^{(1)} - (a'_{13})^{(1)}$ $(S_2)^{(1)} = (a_{15})^{(1)} - (p_{15})^{(1)}$ $(R_1)^{(1)} = (b_{13})^{(1)}(\mu_2)^{(1)} - (b'_{13})^{(1)}$ $(R_2)^{(1)} = (b'_{15})^{(1)} - (r_{15})^{(1)}$ | 290 |
| **Behavior of the solutions of equation** | 291 |
| **Theorem 2:** If we denote and define | |
| **Definition of** $(\sigma_1)^{(2)}, (\sigma_2)^{(2)}, (\tau_1)^{(2)}, (\tau_2)^{(2)}$ : $(\sigma_1)^{(2)}, (\sigma_2)^{(2)}, (\tau_1)^{(2)}, (\tau_2)^{(2)}$ four constants satisfying | 292 |

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| $-(\sigma_2)^{(2)} \le -(a_{16}^{'})^{(2)} + (a_{17}^{'})^{(2)} - (a_{16}^{''})^{(2)}(\mathrm{T}_{17}, t) + (a_{17}^{''})^{(2)}(\mathrm{T}_{17}, t) \le -(\sigma_1)^{(2)}$ | 293 |
| $-(\tau_2)^{(2)} \le -(b_{16}^{'})^{(2)} + (b_{17}^{'})^{(2)} - (b_{16}^{''})^{(2)}\big((G_{19}), t\big) - (b_{17}^{''})^{(2)}\big((G_{19}), t\big) \le -(\tau_1)^{(2)}$ | 294 |
| **Definition of** $(\nu_1)^{(2)}, (\nu_2)^{(2)}, (u_1)^{(2)}, (u_2)^{(2)}$ : | 295 |
| By $(\nu_1)^{(2)} > 0$ , $(\nu_2)^{(2)} < 0$ and respectively $(u_1)^{(2)} > 0$ , $(u_2)^{(2)} < 0$ the roots | 296 |
| of the equations $(a_{17})^{(2)}\big(\nu^{(2)}\big)^2 + (\sigma_1)^{(2)}\nu^{(2)} - (a_{16})^{(2)} = 0$ | 297 |
| and $(b_{14})^{(2)}\big(u^{(2)}\big)^2 + (\tau_1)^{(2)}u^{(2)} - (b_{16})^{(2)} = 0$ and | 298 |
| **Definition of** $(\bar{\nu}_1)^{(2)}, , (\bar{\nu}_2)^{(2)}, (\bar{u}_1)^{(2)}, (\bar{u}_2)^{(2)}$ : | 299 |
| By $(\bar{\nu}_1)^{(2)} > 0$ , $(\bar{\nu}_2)^{(2)} < 0$ and respectively $(\bar{u}_1)^{(2)} > 0$ , $(\bar{u}_2)^{(2)} < 0$ the | 300 |
| roots of the equations $(a_{17})^{(2)}\big(\nu^{(2)}\big)^2 + (\sigma_2)^{(2)}\nu^{(2)} - (a_{16})^{(2)} = 0$ | 301 |
| and $(b_{17})^{(2)}\big(u^{(2)}\big)^2 + (\tau_2)^{(2)}u^{(2)} - (b_{16})^{(2)} = 0$ | 302 |
| **Definition of** $(m_1)^{(2)}$ , $(m_2)^{(2)}$ , $(\mu_1)^{(2)}, (\mu_2)^{(2)}$ :- | 303 |
| If we define $(m_1)^{(2)}$ , $(m_2)^{(2)}$ , $(\mu_1)^{(2)}, (\mu_2)^{(2)}$ by | 304 |
| $(m_2)^{(2)} = (\nu_0)^{(2)}, (m_1)^{(2)} = (\nu_1)^{(2)}$, ***if*** $(\nu_0)^{(2)} < (\nu_1)^{(2)}$ | 305 |
| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ , ***if*** $(\nu_1)^{(2)} < (\nu_0)^{(2)} < (\bar{\nu}_1)^{(2)}$, and $\boxed{(\nu_0)^{(2)} = \frac{\mathrm{G}_{16}^0}{\mathrm{G}_{17}^0}}$ | 306 |
| $(m_2)^{(2)} = (\nu_1)^{(2)}, (m_1)^{(2)} = (\nu_0)^{(2)}$, ***if*** $(\bar{\nu}_1)^{(2)} < (\nu_0)^{(2)}$ | 307 |
| and analogously<br>$(\mu_2)^{(2)} = (u_0)^{(2)}, (\mu_1)^{(2)} = (u_1)^{(2)}$, ***if*** $(u_0)^{(2)} < (u_1)^{(2)}$<br>$(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (\bar{u}_1)^{(2)}$ , ***if*** $(u_1)^{(2)} < (u_0)^{(2)} < (\bar{u}_1)^{(2)}$,<br>and $\boxed{(u_0)^{(2)} = \frac{\mathrm{T}_{16}^0}{\mathrm{T}_{17}^0}}$ | 308 |
| $(\mu_2)^{(2)} = (u_1)^{(2)}, (\mu_1)^{(2)} = (u_0)^{(2)}$, ***if*** $(\bar{u}_1)^{(2)} < (u_0)^{(2)}$ | 309 |
| Then the solution of global equations satisfies the inequalities<br>$\mathrm{G}_{16}^0 \mathrm{e}^{((\mathrm{S}_1)^{(2)} - (p_{16})^{(2)})\mathrm{t}} \le G_{16}(t) \le \mathrm{G}_{16}^0 \mathrm{e}^{(\mathrm{S}_1)^{(2)}\mathrm{t}}$ | 310 |
| $(p_i)^{(2)}$ is defined by equation | |

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| $\frac{1}{(m_1)^{(2)}} \mathrm{G}_{16}^{0} \mathrm{e}^{((\mathrm{S}_1)^{(2)}-(p_{16})^{(2)})\mathrm{t}} \leq G_{17}(t) \leq \frac{1}{(m_2)^{(2)}} \mathrm{G}_{16}^{0} \mathrm{e}^{(\mathrm{S}_1)^{(2)}\mathrm{t}}$ | 311 |
| $\left(\frac{(a_{18})^{(2)}\mathrm{G}_{16}^{0}}{(m_1)^{(2)}((\mathrm{S}_1)^{(2)}-(p_{16})^{(2)}-(\mathrm{S}_2)^{(2)})}\left[\mathrm{e}^{((\mathrm{S}_1)^{(2)}-(p_{16})^{(2)})\mathrm{t}} - \mathrm{e}^{-(\mathrm{S}_2)^{(2)}\mathrm{t}}\right] + \mathrm{G}_{18}^{0}\mathrm{e}^{-(\mathrm{S}_2)^{(2)}\mathrm{t}} \leq \mathrm{G}_{18}(t) \leq \right.$ $\left.\frac{(a_{18})^{(2)}\mathrm{G}_{16}^{0}}{(m_2)^{(2)}((\mathrm{S}_1)^{(2)}-(a_{18}^{'})^{(2)})}\left[\mathrm{e}^{(\mathrm{S}_1)^{(2)}\mathrm{t}} - \mathrm{e}^{-(a_{18}^{'})^{(2)}\mathrm{t}}\right] + \mathrm{G}_{18}^{0}\mathrm{e}^{-(a_{18}^{'})^{(2)}\mathrm{t}}\right)$ | 312 |
| $\boxed{\mathrm{T}_{16}^{0}\mathrm{e}^{(\mathrm{R}_1)^{(2)}t} \leq T_{16}(t) \leq \mathrm{T}_{16}^{0}\mathrm{e}^{((\mathrm{R}_1)^{(2)}+(r_{16})^{(2)})t}}$ | 313 |
| $\frac{1}{(\mu_1)^{(2)}} \mathrm{T}_{16}^{0}\mathrm{e}^{(\mathrm{R}_1)^{(2)}t} \leq T_{16}(t) \leq \frac{1}{(\mu_2)^{(2)}} \mathrm{T}_{16}^{0}\mathrm{e}^{((\mathrm{R}_1)^{(2)}+(r_{16})^{(2)})t}$ | 314 |
| $\frac{(b_{18})^{(2)}\mathrm{T}_{16}^{0}}{(\mu_1)^{(2)}((\mathrm{R}_1)^{(2)}-(b_{18}^{'})^{(2)})}\left[\mathrm{e}^{(\mathrm{R}_1)^{(2)}t} - \mathrm{e}^{-(b_{18}^{'})^{(2)}t}\right] + \mathrm{T}_{18}^{0}\mathrm{e}^{-(b_{18}^{'})^{(2)}t} \leq T_{18}(t) \leq$ $\frac{(a_{18})^{(2)}\mathrm{T}_{16}^{0}}{(\mu_2)^{(2)}((\mathrm{R}_1)^{(2)}+(r_{16})^{(2)}+(\mathrm{R}_2)^{(2)})}\left[\mathrm{e}^{((\mathrm{R}_1)^{(2)}+(r_{16})^{(2)})t} - \mathrm{e}^{-(\mathrm{R}_2)^{(2)}t}\right] + \mathrm{T}_{18}^{0}\mathrm{e}^{-(\mathrm{R}_2)^{(2)}t}$ | 315 |
| **Definition of** $(\mathrm{S}_1)^{(2)}, (\mathrm{S}_2)^{(2)}, (\mathrm{R}_1)^{(2)}, (\mathrm{R}_2)^{(2)}$:- | 316 |
| Where $(\mathrm{S}_1)^{(2)} = (a_{16})^{(2)}(m_2)^{(2)} - (a_{16}^{'})^{(2)}$ $(\mathrm{S}_2)^{(2)} = (a_{18})^{(2)} - (p_{18})^{(2)}$ | 317 |
| $(R_1)^{(2)} = (b_{16})^{(2)}(\mu_2)^{(1)} - (b_{16}^{'})^{(2)}$ $(\mathrm{R}_2)^{(2)} = (b_{18}^{'})^{(2)} - (r_{18})^{(2)}$ | 318 |
| **Behavior of the solutions** **Theorem 3:** If we denote and define **Definition of** $(\sigma_1)^{(3)}, (\sigma_2)^{(3)}, (\tau_1)^{(3)}, (\tau_2)^{(3)}$ : $(\sigma_1)^{(3)}, (\sigma_2)^{(3)}, (\tau_1)^{(3)}, (\tau_2)^{(3)}$ four constants satisfying $-(\sigma_2)^{(3)} \leq -(a_{20}^{'})^{(3)} + (a_{21}^{'})^{(3)} - (a_{20}^{''})^{(3)}(T_{21}, t) + (a_{21}^{''})^{(3)}(T_{21}, t) \leq -(\sigma_1)^{(3)}$ $-(\tau_2)^{(3)} \leq -(b_{20}^{'})^{(3)} + (b_{21}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23}, t) - (b_{21}^{''})^{(3)}((G_{23}), t) \leq -(\tau_1)^{(3)}$ | 319 |
| **Definition of** $(\nu_1)^{(3)}, (\nu_2)^{(3)}, (u_1)^{(3)}, (u_2)^{(3)}$ : By $(\nu_1)^{(3)} > 0$, $(\nu_2)^{(3)} < 0$ and respectively $(u_1)^{(3)} > 0$, $(u_2)^{(3)} < 0$ the roots of the equations $(a_{21})^{(3)}(\nu^{(3)})^2 + (\sigma_1)^{(3)}\nu^{(3)} - (a_{20})^{(3)} = 0$ and $(b_{21})^{(3)}(u^{(3)})^2 + (\tau_1)^{(3)}u^{(3)} - (b_{20})^{(3)} = 0$ and By $(\bar{\nu}_1)^{(3)} > 0$, $(\bar{\nu}_2)^{(3)} < 0$ and respectively $(\bar{u}_1)^{(3)} > 0$, $(\bar{u}_2)^{(3)} < 0$ the | 320 |

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| roots of the equations $(a_{21})^{(3)}\left(v^{(3)}\right)^2 + (\sigma_2)^{(3)}v^{(3)} - (a_{20})^{(3)} = 0$<br><br>and $(b_{21})^{(3)}\left(u^{(3)}\right)^2 + (\tau_2)^{(3)}u^{(3)} - (b_{20})^{(3)} = 0$ | |
| **Definition of** $(m_1)^{(3)}, (m_2)^{(3)}, (\mu_1)^{(3)}, (\mu_2)^{(3)}$ :-<br><br>If we define $(m_1)^{(3)}, (m_2)^{(3)}, (\mu_1)^{(3)}, (\mu_2)^{(3)}$ by<br>$(m_2)^{(3)} = (v_0)^{(3)}, (m_1)^{(3)} = (v_1)^{(3)}$, ***if*** $(v_0)^{(3)} < (v_1)^{(3)}$<br><br>$(m_2)^{(3)} = (v_1)^{(3)}, (m_1)^{(3)} = (\bar{v}_1)^{(3)}$, ***if*** $(v_1)^{(3)} < (v_0)^{(3)} < (\bar{v}_1)^{(3)}$,<br><br>and $\boxed{(v_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0}}$<br><br>$(m_2)^{(3)} = (v_1)^{(3)}, (m_1)^{(3)} = (v_0)^{(3)}$, ***if*** $(\bar{v}_1)^{(3)} < (v_0)^{(3)}$ | 321 |
| and analogously<br><br>$(\mu_2)^{(3)} = (u_0)^{(3)}, (\mu_1)^{(3)} = (u_1)^{(3)}$, ***if*** $(u_0)^{(3)} < (u_1)^{(3)}$<br><br>$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (\bar{u}_1)^{(3)}$, ***if*** $(u_1)^{(3)} < (u_0)^{(3)} < (\bar{u}_1)^{(3)}$, and $\boxed{(u_0)^{(3)} = \frac{T_{20}^0}{T_{21}^0}}$<br><br>$(\mu_2)^{(3)} = (u_1)^{(3)}, (\mu_1)^{(3)} = (u_0)^{(3)}$, ***if*** $(\bar{u}_1)^{(3)} < (u_0)^{(3)}$<br><br>Then the solution of global equations satisfies the inequalities<br><br>$G_{20}^0 e^{((S_1)^{(3)} - (p_{20})^{(3)})t} \le G_{20}(t) \le G_{20}^0 e^{(S_1)^{(3)}t}$<br><br>$(p_i)^{(3)}$ is defined by equation | 322 |
| $\frac{1}{(m_1)^{(3)}} G_{20}^0 e^{((S_1)^{(3)} - (p_{20})^{(3)})t} \le G_{21}(t) \le \frac{1}{(m_2)^{(3)}} G_{20}^0 e^{(S_1)^{(3)}t}$ | 323 |
| $\left(\frac{(a_{22})^{(3)} G_{20}^0}{(m_1)^{(3)}((S_1)^{(3)} - (p_{20})^{(3)} - (S_2)^{(3)})}\left[e^{((S_1)^{(3)} - (p_{20})^{(3)})t} - e^{-(S_2)^{(3)}t}\right] + G_{22}^0 e^{-(S_2)^{(3)}t} \le G_{22}(t) \le \right.$<br>$\left.\frac{(a_{22})^{(3)} G_{20}^0}{(m_2)^{(3)}((S_1)^{(3)} - (a'_{22})^{(3)})}\left[e^{(S_1)^{(3)}t} - e^{-(a'_{22})^{(3)}t}\right] + G_{22}^0 e^{-(a'_{22})^{(3)}t}\right)$ | 324 |
| $\boxed{T_{20}^0 e^{(R_1)^{(3)}t} \le T_{20}(t) \le T_{20}^0 e^{((R_1)^{(3)} + (r_{20})^{(3)})t}}$ | 325 |
| $\frac{1}{(\mu_1)^{(3)}} T_{20}^0 e^{(R_1)^{(3)}t} \le T_{20}(t) \le \frac{1}{(\mu_2)^{(3)}} T_{20}^0 e^{((R_1)^{(3)} + (r_{20})^{(3)})t}$ | 326 |
| $\frac{(b_{22})^{(3)} T_{20}^0}{(\mu_1)^{(3)}((R_1)^{(3)} - (b'_{22})^{(3)})}\left[e^{(R_1)^{(3)}t} - e^{-(b'_{22})^{(3)}t}\right] + T_{22}^0 e^{-(b'_{22})^{(3)}t} \le T_{22}(t) \le$<br><br>$\frac{(a_{22})^{(3)} T_{20}^0}{(\mu_2)^{(3)}((R_1)^{(3)} + (r_{20})^{(3)} + (R_2)^{(3)})}\left[e^{((R_1)^{(3)} + (r_{20})^{(3)})t} - e^{-(R_2)^{(3)}t}\right] + T_{22}^0 e^{-(R_2)^{(3)}t}$ | 327 |

**Definition of** $(S_1)^{(3)}, (S_2)^{(3)}, (R_1)^{(3)}, (R_2)^{(3)}$:- 328

Where $(S_1)^{(3)} = (a_{20})^{(3)}(m_2)^{(3)} - (a'_{20})^{(3)}$

$(S_2)^{(3)} = (a_{22})^{(3)} - (p_{22})^{(3)}$

$(R_1)^{(3)} = (b_{20})^{(3)}(\mu_2)^{(3)} - (b'_{20})^{(3)}$

$(R_2)^{(3)} = (b'_{22})^{(3)} - (r_{22})^{(3)}$

**Behavior of the solutions of equation**

**Theorem:** If we denote and define

**Definition of** $(\sigma_1)^{(4)}, (\sigma_2)^{(4)}, (\tau_1)^{(4)}, (\tau_2)^{(4)}$ :

$(\sigma_1)^{(4)}, (\sigma_2)^{(4)}, (\tau_1)^{(4)}, (\tau_2)^{(4)}$ four constants satisfying

$-(\sigma_2)^{(4)} \le -(a'_{24})^{(4)} + (a'_{25})^{(4)} - (a''_{24})^{(4)}(T_{25}, t) + (a''_{25})^{(4)}(T_{25}, t) \le -(\sigma_1)^{(4)}$

$-(\tau_2)^{(4)} \le -(b'_{24})^{(4)} + (b'_{25})^{(4)} - (b''_{24})^{(4)}((G_{27}), t) - (b''_{25})^{(4)}((G_{27}), t) \le -(\tau_1)^{(4)}$

**Definition of** $(v_1)^{(4)}, (v_2)^{(4)}, (u_1)^{(4)}, (u_2)^{(4)}, v^{(4)}, u^{(4)}$ : 329

By $(v_1)^{(4)} > 0$, $(v_2)^{(4)} < 0$ and respectively $(u_1)^{(4)} > 0$, $(u_2)^{(4)} < 0$ the roots of the equations $(a_{25})^{(4)}\left(v^{(4)}\right)^2 + (\sigma_1)^{(4)}v^{(4)} - (a_{24})^{(4)} = 0$
and $(b_{25})^{(4)}\left(u^{(4)}\right)^2 + (\tau_1)^{(4)}u^{(4)} - (b_{24})^{(4)} = 0$ and

**Definition of** $(\bar{v}_1)^{(4)}, , (\bar{v}_2)^{(4)}, (\bar{u}_1)^{(4)}, (\bar{u}_2)^{(4)}$ : 330

By $(\bar{v}_1)^{(4)} > 0$, $(\bar{v}_2)^{(4)} < 0$ and respectively $(\bar{u}_1)^{(4)} > 0$, $(\bar{u}_2)^{(4)} < 0$ the roots of the equations $(a_{25})^{(4)}\left(v^{(4)}\right)^2 + (\sigma_2)^{(4)}v^{(4)} - (a_{24})^{(4)} = 0$
and $(b_{25})^{(4)}\left(u^{(4)}\right)^2 + (\tau_2)^{(4)}u^{(4)} - (b_{24})^{(4)} = 0$

**Definition of** $(m_1)^{(4)}, (m_2)^{(4)}, (\mu_1)^{(4)}, (\mu_2)^{(4)}, (v_0)^{(4)}$ :-

If we define $(m_1)^{(4)}, (m_2)^{(4)}, (\mu_1)^{(4)}, (\mu_2)^{(4)}$ by

$(m_2)^{(4)} = (v_0)^{(4)}, (m_1)^{(4)} = (v_1)^{(4)}$, ***if*** $(v_0)^{(4)} < (v_1)^{(4)}$

$(m_2)^{(4)} = (v_1)^{(4)}, (m_1)^{(4)} = (\bar{v}_1)^{(4)}$, ***if*** $(v_4)^{(4)} < (v_0)^{(4)} < (\bar{v}_1)^{(4)}$,
and $\boxed{(v_0)^{(4)} = \frac{G^0_{24}}{G^0_{25}}}$

$(m_2)^{(4)} = (v_4)^{(4)}, (m_1)^{(4)} = (v_0)^{(4)}$, ***if*** $(\bar{v}_4)^{(4)} < (v_0)^{(4)}$

and analogously 331

$(\mu_2)^{(4)} = (u_0)^{(4)}, (\mu_1)^{(4)} = (u_1)^{(4)}$, ***if*** $(u_0)^{(4)} < (u_1)^{(4)}$

$(\mu_2)^{(4)} = (u_1)^{(4)}, (\mu_1)^{(4)} = (\bar{u}_1)^{(4)}$, ***if*** $(u_1)^{(4)} < (u_0)^{(4)} < (\bar{u}_1)^{(4)}$,

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| and $\boxed{(u_0)^{(4)} = \frac{T_{24}^0}{T_{25}^0}}$<br>$(\mu_2)^{(4)} = (u_1)^{(4)}, (\mu_1)^{(4)} = (u_0)^{(4)}, \boldsymbol{if}\ (\bar{u}_1)^{(4)} < (u_0)^{(4)}$ where $(u_1)^{(4)}, (\bar{u}_1)^{(4)}$ | |
| Then the solution of global equations satisfies the inequalities<br>$G_{24}^0 e^{((S_1)^{(4)} - (p_{24})^{(4)})t} \le G_{24}(t) \le G_{24}^0 e^{(S_1)^{(4)} t}$<br>where $(p_i)^{(4)}$ is defined by equation | 332 |
| $\frac{1}{(m_1)^{(4)}} G_{24}^0 e^{((S_1)^{(4)} - (p_{24})^{(4)})t} \le G_{25}(t) \le \frac{1}{(m_2)^{(4)}} G_{24}^0 e^{(S_1)^{(4)} t}$ | 333 |
| $\left(\frac{(a_{26})^{(4)} G_{24}^0}{(m_1)^{(4)}((S_1)^{(4)} - (p_{24})^{(4)} - (S_2)^{(4)})}\right)\left[e^{((S_1)^{(4)} - (p_{24})^{(4)})t} - e^{-(S_2)^{(4)} t}\right] + G_{26}^0 e^{-(S_2)^{(4)} t} \le G_{26}(t) \le$<br>*(a26)4G240(m2)4(S1)4−(a26′)4e(S1)4t−e−(a26′)4t+ G260e−(a26′)4t* | 334 |
| $\boxed{T_{24}^0 e^{(R_1)^{(4)} t} \le T_{24}(t) \le T_{24}^0 e^{((R_1)^{(4)} + (r_{24})^{(4)})t}}$ | |
| $\frac{1}{(\mu_1)^{(4)}} T_{24}^0 e^{(R_1)^{(4)} t} \le T_{24}(t) \le \frac{1}{(\mu_2)^{(4)}} T_{24}^0 e^{((R_1)^{(4)} + (r_{24})^{(4)})t}$ | 335 |
| $\frac{(b_{26})^{(4)} T_{24}^0}{(\mu_1)^{(4)}((R_1)^{(4)} - (b_{26}')^{(4)})}\left[e^{(R_1)^{(4)} t} - e^{-(b_{26}')^{(4)} t}\right] + T_{26}^0 e^{-(b_{26}')^{(4)} t} \le T_{26}(t) \le$<br>$\frac{(a_{26})^{(4)} T_{24}^0}{(\mu_2)^{(4)}((R_1)^{(4)} + (r_{24})^{(4)} + (R_2)^{(4)})}\left[e^{((R_1)^{(4)} + (r_{24})^{(4)})t} - e^{-(R_2)^{(4)} t}\right] + T_{26}^0 e^{-(R_2)^{(4)} t}$ | 336 |
| **Definition of** $(S_1)^{(4)}, (S_2)^{(4)}, (R_1)^{(4)}, (R_2)^{(4)}$:-<br>Where $(S_1)^{(4)} = (a_{24})^{(4)}(m_2)^{(4)} - (a_{24}')^{(4)}$<br>$(S_2)^{(4)} = (a_{26})^{(4)} - (p_{26})^{(4)}$<br>$(R_1)^{(4)} = (b_{24})^{(4)}(\mu_2)^{(4)} - (b_{24}')^{(4)}$<br>$(R_2)^{(4)} = (b_{26}')^{(4)} - (r_{26})^{(4)}$ | 337 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ :<br>$(\sigma_1)^{(5)}, (\sigma_2)^{(5)}, (\tau_1)^{(5)}, (\tau_2)^{(5)}$ four constants satisfying<br>$-(\sigma_2)^{(5)} \le -(a_{28}')^{(5)} + (a_{29}')^{(5)} - (a_{28}'')^{(5)}(T_{29}, t) + (a_{29}'')^{(5)}(T_{29}, t) \le -(\sigma_1)^{(5)}$<br>$-(\tau_2)^{(5)} \le -(b_{28}')^{(5)} + (b_{29}')^{(5)} - (b_{28}'')^{(5)}((G_{31}), t) - (b_{29}'')^{(5)}((G_{31}), t) \le -(\tau_1)^{(5)}$ | 338 |

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| **Definition of** $(v_1)^{(5)}, (v_2)^{(5)}, (u_1)^{(5)}, (u_2)^{(5)}, v^{(5)}, u^{(5)}$ :<br><br>By $(v_1)^{(5)} > 0$ , $(v_2)^{(5)} < 0$ and respectively $(u_1)^{(5)} > 0$ , $(u_2)^{(5)} < 0$ the roots of the equations $(a_{29})^{(5)}\left(v^{(5)}\right)^2 + (\sigma_1)^{(5)}v^{(5)} - (a_{28})^{(5)} = 0$<br>and $(b_{29})^{(5)}\left(u^{(5)}\right)^2 + (\tau_1)^{(5)}u^{(5)} - (b_{28})^{(5)} = 0$ and | 339 |
| **Definition of** $(\bar{v}_1)^{(5)},, (\bar{v}_2)^{(5)}, (\bar{u}_1)^{(5)}, (\bar{u}_2)^{(5)}$ :<br><br>By $(\bar{v}_1)^{(5)} > 0$ , $(\bar{v}_2)^{(5)} < 0$ and respectively $(\bar{u}_1)^{(5)} > 0$ , $(\bar{u}_2)^{(5)} < 0$ the roots of the equations $(a_{29})^{(5)}\left(v^{(5)}\right)^2 + (\sigma_2)^{(5)}v^{(5)} - (a_{28})^{(5)} = 0$<br>and $(b_{29})^{(5)}\left(u^{(5)}\right)^2 + (\tau_2)^{(5)}u^{(5)} - (b_{28})^{(5)} = 0$<br>**Definition of** $(m_1)^{(5)}$ , $(m_2)^{(5)}$ , $(\mu_1)^{(5)}, (\mu_2)^{(5)}, (v_0)^{(5)}$ :-<br><br>If we define $(m_1)^{(5)}$ , $(m_2)^{(5)}$ , $(\mu_1)^{(5)}, (\mu_2)^{(5)}$ by<br><br>$(m_2)^{(5)} = (v_0)^{(5)}, (m_1)^{(5)} = (v_1)^{(5)}$, ***if*** $(v_0)^{(5)} < (v_1)^{(5)}$<br><br>$(m_2)^{(5)} = (v_1)^{(5)}, (m_1)^{(5)} = (\bar{v}_1)^{(5)}$ , ***if*** $(v_1)^{(5)} < (v_0)^{(5)} < (\bar{v}_1)^{(5)}$,<br>and $\boxed{(v_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}}$<br><br>$(m_2)^{(5)} = (v_1)^{(5)}, (m_1)^{(5)} = (v_0)^{(5)}$, ***if*** $(\bar{v}_1)^{(5)} < (v_0)^{(5)}$ | 340 |
| and analogously<br><br>$(\mu_2)^{(5)} = (u_0)^{(5)}, (\mu_1)^{(5)} = (u_1)^{(5)}$, ***if*** $(u_0)^{(5)} < (u_1)^{(5)}$<br><br>$(\mu_2)^{(5)} = (u_1)^{(5)}, (\mu_1)^{(5)} = (\bar{u}_1)^{(5)}$ , ***if*** $(u_1)^{(5)} < (u_0)^{(5)} < (\bar{u}_1)^{(5)}$,<br>and $\boxed{(u_0)^{(5)} = \frac{T_{28}^0}{T_{29}^0}}$<br><br>$(\mu_2)^{(5)} = (u_1)^{(5)}, (\mu_1)^{(5)} = (u_0)^{(5)}$, ***if*** $(\bar{u}_1)^{(5)} < (u_0)^{(5)}$ where $(u_1)^{(5)}, (\bar{u}_1)^{(5)}$ | 341 |
| Then the solution of global equations satisfies the inequalities<br><br>$G_{28}^0 e^{((S_1)^{(5)} - (p_{28})^{(5)})t} \le G_{28}(t) \le G_{28}^0 e^{(S_1)^{(5)}t}$<br><br>where $(p_i)^{(5)}$ is defined by equation | 342 |
| $\frac{1}{(m_5)^{(5)}} G_{28}^0 e^{((S_1)^{(5)} - (p_{28})^{(5)})t} \le G_{29}(t) \le \frac{1}{(m_2)^{(5)}} G_{28}^0 e^{(S_1)^{(5)}t}$ | 343 |
| $\left(\frac{(a_{30})^{(5)} G_{28}^0}{(m_1)^{(5)}((S_1)^{(5)} - (p_{28})^{(5)} - (S_2)^{(5)})}\left[e^{((S_1)^{(5)} - (p_{28})^{(5)})t} - e^{-(S_2)^{(5)}t}\right] + G_{30}^0 e^{-(S_2)^{(5)}t} \le G_{30}(t) \le$<br>*(a30)5G280(m2)5(S1)5−(a30′)5e(S1)5t−e−(a30′)5t+ G300e−(a30′)5t* | 344 |
| $\boxed{T_{28}^0 e^{(R_1)^{(5)}t} \le T_{28}(t) \le T_{28}^0 e^{((R_1)^{(5)} + (r_{28})^{(5)})t}}$ | 345 |

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| $\frac{1}{(\mu_1)^{(5)}}T_{28}^0 e^{(R_1)^{(5)}t} \le T_{28}(t) \le \frac{1}{(\mu_2)^{(5)}}T_{28}^0 e^{((R_1)^{(5)}+(r_{28})^{(5)})t}$ | 346 |
| $\frac{(b_{30})^{(5)}T_{28}^0}{(\mu_1)^{(5)}((R_1)^{(5)}-(b_{30}')^{(5)})}\left[e^{(R_1)^{(5)}t} - e^{-(b_{30}')^{(5)}t}\right] + T_{30}^0 e^{-(b_{30}')^{(5)}t} \le T_{30}(t) \le$<br>$\frac{(a_{30})^{(5)}T_{28}^0}{(\mu_2)^{(5)}((R_1)^{(5)}+(r_{28})^{(5)}+(R_2)^{(5)})}\left[e^{((R_1)^{(5)}+(r_{28})^{(5)})t} - e^{-(R_2)^{(5)}t}\right] + T_{30}^0 e^{-(R_2)^{(5)}t}$ | 347 |
| **Definition of** $(S_1)^{(5)}, (S_2)^{(5)}, (R_1)^{(5)}, (R_2)^{(5)}$:-<br>Where $(S_1)^{(5)} = (a_{28})^{(5)}(m_2)^{(5)} - (a_{28}')^{(5)}$<br>$(S_2)^{(5)} = (a_{30})^{(5)} - (p_{30})^{(5)}$<br>$(R_1)^{(5)} = (b_{28})^{(5)}(\mu_2)^{(5)} - (b_{28}')^{(5)}$<br>$(R_2)^{(5)} = (b_{30}')^{(5)} - (r_{30})^{(5)}$ | 348 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(6)}, (\sigma_2)^{(6)}, (\tau_1)^{(6)}, (\tau_2)^{(6)}$ :<br>$(\sigma_1)^{(6)}, (\sigma_2)^{(6)}, (\tau_1)^{(6)}, (\tau_2)^{(6)}$ four constants satisfying<br>$-(\sigma_2)^{(6)} \le -(a_{32}')^{(6)} + (a_{33}')^{(6)} - (a_{32}'')^{(6)}(T_{33}, t) + (a_{33}'')^{(6)}(T_{33}, t) \le -(\sigma_1)^{(6)}$<br>$-(\tau_2)^{(6)} \le -(b_{32}')^{(6)} + (b_{33}')^{(6)} - (b_{32}'')^{(6)}((G_{35}), t) - (b_{33}'')^{(6)}((G_{35}), t) \le -(\tau_1)^{(6)}$ | 349 |
| **Definition of** $(v_1)^{(6)}, (v_2)^{(6)}, (u_1)^{(6)}, (u_2)^{(6)}, v^{(6)}, u^{(6)}$ :<br>By $(v_1)^{(6)} > 0$, $(v_2)^{(6)} < 0$ and respectively $(u_1)^{(6)} > 0$, $(u_2)^{(6)} < 0$ the roots of the equations $(a_{33})^{(6)}(v^{(6)})^2 + (\sigma_1)^{(6)}v^{(6)} - (a_{32})^{(6)} = 0$<br>and $(b_{33})^{(6)}(u^{(6)})^2 + (\tau_1)^{(6)}u^{(6)} - (b_{32})^{(6)} = 0$ and | 350 |
| **Definition of** $(\bar{v}_1)^{(6)},, (\bar{v}_2)^{(6)}, (\bar{u}_1)^{(6)}, (\bar{u}_2)^{(6)}$ :<br>By $(\bar{v}_1)^{(6)} > 0$, $(\bar{v}_2)^{(6)} < 0$ and respectively $(\bar{u}_1)^{(6)} > 0$, $(\bar{u}_2)^{(6)} < 0$ the roots of the equations $(a_{33})^{(6)}(v^{(6)})^2 + (\sigma_2)^{(6)}v^{(6)} - (a_{32})^{(6)} = 0$<br>and $(b_{33})^{(6)}(u^{(6)})^2 + (\tau_2)^{(6)}u^{(6)} - (b_{32})^{(6)} = 0$<br>**Definition of** $(m_1)^{(6)}, (m_2)^{(6)}, (\mu_1)^{(6)}, (\mu_2)^{(6)}, (v_0)^{(6)}$ :-<br>If we define $(m_1)^{(6)}, (m_2)^{(6)}, (\mu_1)^{(6)}, (\mu_2)^{(6)}$ by<br>$(m_2)^{(6)} = (v_0)^{(6)}, (m_1)^{(6)} = (v_1)^{(6)}$, ***if*** $(v_0)^{(6)} < (v_1)^{(6)}$<br>$(m_2)^{(6)} = (v_1)^{(6)}, (m_1)^{(6)} = (\bar{v}_6)^{(6)}$, ***if*** $(v_1)^{(6)} < (v_0)^{(6)} < (\bar{v}_1)^{(6)}$, | 351 |

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| and $\boxed{(v_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0}}$<br>$(m_2)^{(6)} = (v_1)^{(6)}, (m_1)^{(6)} = (v_0)^{(6)},$ ***if*** $(\bar{v}_1)^{(6)} < (v_0)^{(6)}$ | |
| and analogously<br>$(\mu_2)^{(6)} = (u_0)^{(6)}, (\mu_1)^{(6)} = (u_1)^{(6)},$ ***if*** $(u_0)^{(6)} < (u_1)^{(6)}$<br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (\bar{u}_1)^{(6)}$ , ***if*** $(u_1)^{(6)} < (u_0)^{(6)} < (\bar{u}_1)^{(6)}$,<br>and $\boxed{(u_0)^{(6)} = \frac{T_{32}^0}{T_{33}^0}}$<br>$(\mu_2)^{(6)} = (u_1)^{(6)}, (\mu_1)^{(6)} = (u_0)^{(6)},$ ***if*** $(\bar{u}_1)^{(6)} < (u_0)^{(6)}$ where $(u_1)^{(6)}, (\bar{u}_1)^{(6)}$ | 352 |
| Then the solution of global equations satisfies the inequalities<br>$G_{32}^0 e^{((S_1)^{(6)} - (p_{32})^{(6)})t} \le G_{32}(t) \le G_{32}^0 e^{(S_1)^{(6)} t}$<br>where $(p_i)^{(6)}$ is defined by equation | 353 |
| $\frac{1}{(m_1)^{(6)}} G_{32}^0 e^{((S_1)^{(6)} - (p_{32})^{(6)})t} \le G_{33}(t) \le \frac{1}{(m_2)^{(6)}} G_{32}^0 e^{(S_1)^{(6)} t}$ | 354 |
| $\left(\frac{(a_{34})^{(6)} G_{32}^0}{(m_1)^{(6)}((S_1)^{(6)} - (p_{32})^{(6)} - (S_2)^{(6)})}\right)\left[e^{((S_1)^{(6)} - (p_{32})^{(6)})t} - e^{-(S_2)^{(6)} t}\right] + G_{34}^0 e^{-(S_2)^{(6)} t} \le G_{34}(t) \le$<br>*(a34)6G320(m2)6(S1)6−(a34')6e(S1)6t−e−(a34')6t+ G340e−(a34')6t* | 355 |
| $\boxed{T_{32}^0 e^{(R_1)^{(6)} t} \le T_{32}(t) \le T_{32}^0 e^{((R_1)^{(6)} + (r_{32})^{(6)})t}}$ | 356 |
| $\frac{1}{(\mu_1)^{(6)}} T_{32}^0 e^{(R_1)^{(6)} t} \le T_{32}(t) \le \frac{1}{(\mu_2)^{(6)}} T_{32}^0 e^{((R_1)^{(6)} + (r_{32})^{(6)})t}$ | 357 |
| $\frac{(b_{34})^{(6)} T_{32}^0}{(\mu_1)^{(6)}((R_1)^{(6)} - (b_{34}')^{(6)})}\left[e^{(R_1)^{(6)} t} - e^{-(b_{34}')^{(6)} t}\right] + T_{34}^0 e^{-(b_{34}')^{(6)} t} \le T_{34}(t) \le$<br>$\frac{(a_{34})^{(6)} T_{32}^0}{(\mu_2)^{(6)}((R_1)^{(6)} + (r_{32})^{(6)} + (R_2)^{(6)})}\left[e^{((R_1)^{(6)} + (r_{32})^{(6)})t} - e^{-(R_2)^{(6)} t}\right] + T_{34}^0 e^{-(R_2)^{(6)} t}$ | 358 |
| **Definition of** $(S_1)^{(6)}, (S_2)^{(6)}, (R_1)^{(6)}, (R_2)^{(6)}$:-<br>Where $(S_1)^{(6)} = (a_{32})^{(6)}(m_2)^{(6)} - (a_{32}')^{(6)}$<br>$(S_2)^{(6)} = (a_{34})^{(6)} - (p_{34})^{(6)}$<br>$(R_1)^{(6)} = (b_{32})^{(6)}(\mu_2)^{(6)} - (b_{32}')^{(6)}$<br>$(R_2)^{(6)} = (b_{34}')^{(6)} - (r_{34})^{(6)}$ | 359 |
| **Behavior of the solutions of equation** | |

**Theorem 2:** If we denote and define

**Definition of** $(\sigma_1)^{(7)}, (\sigma_2)^{(7)}, (\tau_1)^{(7)}, (\tau_2)^{(7)}$ :

$(\sigma_1)^{(7)}, (\sigma_2)^{(7)}, (\tau_1)^{(7)}, (\tau_2)^{(7)}$ four constants satisfying

$$-(\sigma_2)^{(7)} \le -(a_{36}^{'})^{(7)} + (a_{37}^{'})^{(7)} - (a_{36}^{''})^{(7)}(T_{37}, t) + (a_{37}^{''})^{(7)}(T_{37}, t) \le -(\sigma_1)^{(7)}$$

$$-(\tau_2)^{(7)} \le -(b_{36}^{'})^{(7)} + (b_{37}^{'})^{(7)} - (b_{36}^{''})^{(7)}\big((G_{39}), t\big) - (b_{37}^{''})^{(7)}\big((G_{39}), t\big) \le -(\tau_1)^{(7)}$$

**Definition of** $(\nu_1)^{(7)}, (\nu_2)^{(7)}, (u_1)^{(7)}, (u_2)^{(7)}, \nu^{(7)}, u^{(7)}$ : 361

By $(\nu_1)^{(7)} > 0, (\nu_2)^{(7)} < 0$ and respectively $(u_1)^{(7)} > 0, (u_2)^{(7)} < 0$ the roots of the equations $(a_{37})^{(7)}\big(\nu^{(7)}\big)^2 + (\sigma_1)^{(7)}\nu^{(7)} - (a_{36})^{(7)} = 0$
and $(b_{37})^{(7)}\big(u^{(7)}\big)^2 + (\tau_1)^{(7)}u^{(7)} - (b_{36})^{(7)} = 0$ and

**Definition of** $(\bar{\nu}_1)^{(7)}, , (\bar{\nu}_2)^{(7)}, (\bar{u}_1)^{(7)}, (\bar{u}_2)^{(7)}$ : 362

By $(\bar{\nu}_1)^{(7)} > 0, (\bar{\nu}_2)^{(7)} < 0$ and respectively $(\bar{u}_1)^{(7)} > 0, (\bar{u}_2)^{(7)} < 0$ the roots of the equations $(a_{37})^{(7)}\big(\nu^{(7)}\big)^2 + (\sigma_2)^{(7)}\nu^{(7)} - (a_{36})^{(7)} = 0$
and $(b_{37})^{(7)}\big(u^{(7)}\big)^2 + (\tau_2)^{(7)}u^{(7)} - (b_{36})^{(7)} = 0$

**Definition of** $(m_1)^{(7)}, (m_2)^{(7)}, (\mu_1)^{(7)}, (\mu_2)^{(7)}, (\nu_0)^{(7)}$ :-

If we define $(m_1)^{(7)}, (m_2)^{(7)}, (\mu_1)^{(7)}, (\mu_2)^{(7)}$ by

$(m_2)^{(7)} = (\nu_0)^{(7)}, (m_1)^{(7)} = (\nu_1)^{(7)}$, ***if*** $(\nu_0)^{(7)} < (\nu_1)^{(7)}$

$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\bar{\nu}_1)^{(7)}$, ***if*** $(\nu_1)^{(7)} < (\nu_0)^{(7)} < (\bar{\nu}_1)^{(7)}$,
and $\boxed{(\nu_0)^{(7)} = \frac{G_{36}^0}{G_{37}^0}}$

$(m_2)^{(7)} = (\nu_1)^{(7)}, (m_1)^{(7)} = (\nu_0)^{(7)}$, ***if*** $(\bar{\nu}_1)^{(7)} < (\nu_0)^{(7)}$

and analogously 363

$(\mu_2)^{(7)} = (u_0)^{(7)}, (\mu_1)^{(7)} = (u_1)^{(7)}$, ***if*** $(u_0)^{(7)} < (u_1)^{(7)}$

$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (\bar{u}_1)^{(7)}$, ***if*** $(u_1)^{(7)} < (u_0)^{(7)} < (\bar{u}_1)^{(7)}$,
and $\boxed{(u_0)^{(7)} = \frac{T_{36}^0}{T_{37}^0}}$

$(\mu_2)^{(7)} = (u_1)^{(7)}, (\mu_1)^{(7)} = (u_0)^{(7)}$, ***if*** $(\bar{u}_1)^{(7)} < (u_0)^{(7)}$ where $(u_1)^{(7)}, (\bar{u}_1)^{(7)}$

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| Then the solution of global equations satisfies the inequalities<br>$G_{36}^0 e^{((S_1)^{(7)}-(p_{36})^{(7)})t} \le G_{36}(t) \le G_{36}^0 e^{(S_1)^{(7)}t}$<br>where $(p_i)^{(7)}$ is defined by equation | 364 |
| $\frac{1}{(m_7)^{(7)}} G_{36}^0 e^{((S_1)^{(7)}-(p_{36})^{(7)})t} \le G_{37}(t) \le \frac{1}{(m_2)^{(7)}} G_{36}^0 e^{(S_1)^{(7)}t}$ | 365 |
| $(\frac{(a_{38})^{(7)} G_{36}^0}{(m_1)^{(7)}((S_1)^{(7)}-(p_{36})^{(7)}-(S_2)^{(7)})}\left[e^{((S_1)^{(7)}-(p_{36})^{(7)})t} - e^{-(S_2)^{(7)}t}\right] + G_{38}^0 e^{-(S_2)^{(7)}t} \le G_{38}(t) \le$<br>$\frac{(a_{38})^{(7)} G_{36}^0}{(m_2)^{(7)}((S_1)^{(7)}-(a_{38}')^{(7)})}[e^{(S_1)^{(7)}t} - e^{-(a_{38}')^{(7)}t}] + G_{38}^0 e^{-(a_{38}')^{(7)}t})$ | 366 |
| $\boxed{T_{36}^0 e^{(R_1)^{(7)}t} \le T_{36}(t) \le T_{36}^0 e^{((R_1)^{(7)}+(r_{36})^{(7)})t}}$ | 367 |
| $\frac{1}{(\mu_1)^{(7)}} T_{36}^0 e^{(R_1)^{(7)}t} \le T_{36}(t) \le \frac{1}{(\mu_2)^{(7)}} T_{36}^0 e^{((R_1)^{(7)}+(r_{36})^{(7)})t}$ | 368 |
| $\frac{(b_{38})^{(7)} T_{36}^0}{(\mu_1)^{(7)}((R_1)^{(7)}-(b_{38}')^{(7)})}\left[e^{(R_1)^{(7)}t} - e^{-(b_{38}')^{(7)}t}\right] + T_{38}^0 e^{-(b_{38}')^{(7)}t} \le T_{38}(t) \le$<br>$\frac{(a_{38})^{(7)} T_{36}^0}{(\mu_2)^{(7)}((R_1)^{(7)}+(r_{36})^{(7)}+(R_2)^{(7)})}\left[e^{((R_1)^{(7)}+(r_{36})^{(7)})t} - e^{-(R_2)^{(7)}t}\right] + T_{38}^0 e^{-(R_2)^{(7)}t}$ | 369 |
| **Definition of** $(S_1)^{(7)}, (S_2)^{(7)}, (R_1)^{(7)}, (R_2)^{(7)}$:-<br>Where $(S_1)^{(7)} = (a_{36})^{(7)}(m_2)^{(7)} - (a_{36}')^{(7)}$<br>$(S_2)^{(7)} = (a_{38})^{(7)} - (p_{38})^{(7)}$<br>$(R_1)^{(7)} = (b_{36})^{(7)}(\mu_2)^{(7)} - (b_{36}')^{(7)}$<br>$(R_2)^{(7)} = (b_{38}')^{(7)} - (r_{38})^{(7)}$ | 370 |
| **Behavior of the solutions of equation**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ :<br>$(\sigma_1)^{(8)}, (\sigma_2)^{(8)}, (\tau_1)^{(8)}, (\tau_2)^{(8)}$ four constants satisfying<br>$-(\sigma_2)^{(8)} \le -(a_{40}')^{(8)} + (a_{41}')^{(8)} - (a_{40}'')^{(8)}(T_{41}, t) + (a_{41}'')^{(8)}(T_{41}, t) \le -(\sigma_1)^{(8)}$<br>$-(\tau_2)^{(8)} \le -(b_{40}')^{(8)} + (b_{41}')^{(8)} - (b_{40}'')^{(8)}((G_{43}), t) - (b_{41}'')^{(8)}((G_{43}), t) \le -(\tau_1)^{(8)}$ | 371 |

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|---|---|
| **Definition of** $(v_1)^{(8)}, (v_2)^{(8)}, (u_1)^{(8)}, (u_2)^{(8)}, v^{(8)}, u^{(8)}$ : <br><br> By $(v_1)^{(8)} > 0$ , $(v_2)^{(8)} < 0$ and respectively $(u_1)^{(8)} > 0$ , $(u_2)^{(8)} < 0$ the roots of the equations $(a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_1)^{(8)}v^{(8)} - (a_{40})^{(8)} = 0$ <br> and $(b_{41})^{(8)}\left(u^{(8)}\right)^2 + (\tau_1)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$ and | 372 |
| **Definition of** $(\bar{v}_1)^{(8)}$,, $(\bar{v}_2)^{(8)}, (\bar{u}_1)^{(8)}, (\bar{u}_2)^{(8)}$ : <br><br> By $(\bar{v}_1)^{(8)} > 0$ , $(\bar{v}_2)^{(8)} < 0$ and respectively $(\bar{u}_1)^{(8)} > 0$ , $(\bar{u}_2)^{(8)} < 0$ the <br><br> roots of the equations $(a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_2)^{(8)}v^{(8)} - (a_{40})^{(8)} = 0$ <br><br> and $(b_{41})^{(8)}\left(u^{(8)}\right)^2 + (\tau_2)^{(8)}u^{(8)} - (b_{40})^{(8)} = 0$ <br><br> **Definition of** $(m_1)^{(8)}$ , $(m_2)^{(8)}$ , $(\mu_1)^{(8)}, (\mu_2)^{(8)}, (v_0)^{(8)}$ :- <br><br> If we define $(m_1)^{(8)}$ , $(m_2)^{(8)}$ , $(\mu_1)^{(8)}, (\mu_2)^{(8)}$ by <br><br> $(m_2)^{(8)} = (v_0)^{(8)}, (m_1)^{(8)} = (v_1)^{(8)}$, ***if*** $(v_0)^{(8)} < (v_1)^{(8)}$ <br><br> $(m_2)^{(8)} = (v_1)^{(8)}, (m_1)^{(8)} = (\bar{v}_1)^{(8)}$ , ***if*** $(v_1)^{(8)} < (v_0)^{(8)} < (\bar{v}_1)^{(8)}$, <br> and $\boxed{(v_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}}$ <br><br> $(m_2)^{(8)} = (v_1)^{(8)}, (m_1)^{(8)} = (v_0)^{(8)}$, ***if*** $(\bar{v}_1)^{(8)} < (v_0)^{(8)}$ | |
| and analogously <br><br> $(\mu_2)^{(8)} = (u_0)^{(8)}, (\mu_1)^{(8)} = (u_1)^{(8)}$, ***if*** $(u_0)^{(8)} < (u_1)^{(8)}$ <br><br> $(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (\bar{u}_1)^{(8)}$ , ***if*** $(u_1)^{(8)} < (u_0)^{(8)} < (\bar{u}_1)^{(8)}$, <br> and $\boxed{(u_0)^{(8)} = \frac{T_{40}^0}{T_{41}^0}}$ <br><br> $(\mu_2)^{(8)} = (u_1)^{(8)}, (\mu_1)^{(8)} = (u_0)^{(8)}$, ***if*** $(\bar{u}_1)^{(8)} < (u_0)^{(8)}$ where $(u_1)^{(8)}, (\bar{u}_1)^{(8)}$ | 374 |
| Then the solution of global equations satisfies the inequalities <br><br> $G_{40}^0 e^{((S_1)^{(8)} - (p_{40})^{(8)})t} \leq G_{40}(t) \leq G_{40}^0 e^{(S_1)^{(8)}t}$ <br><br> where $(p_i)^{(8)}$ is defined by equation | 375 |

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| $\frac{1}{(m_1)^{(8)}} G_{40}^0 e^{((S_1)^{(8)}-(p_{40})^{(8)})t} \le G_{41}(t) \le \frac{1}{(m_2)^{(8)}} G_{40}^0 e^{(S_1)^{(8)}t}$ | 376 |
| $\left(\frac{(a_{42})^{(8)}G_{40}^0}{(m_1)^{(8)}((S_1)^{(8)}-(p_{40})^{(8)}-(S_2)^{(8)})}\left[e^{((S_1)^{(8)}-(p_{40})^{(8)})t} - e^{-(S_2)^{(8)}t}\right] + G_{42}^0 e^{-(S_2)^{(8)}t} \le G_{42}(t) \le \right.$ $\left.\frac{(a_{42})^{(8)}G_{40}^0}{(m_2)^{(8)}((S_1)^{(8)}-(a'_{42})^{(8)})}[e^{(S_1)^{(8)}t} - e^{-(a'_{42})^{(8)}t}] + G_{42}^0 e^{-(a'_{42})^{(8)}t}\right)$ | 377 |
| $\boxed{T_{40}^0 e^{(R_1)^{(8)}t} \le T_{40}(t) \le T_{40}^0 e^{((R_1)^{(8)}+(r_{40})^{(8)})t}}$ | 378 |
| $\frac{1}{(\mu_1)^{(8)}} T_{40}^0 e^{(R_1)^{(8)}t} \le T_{40}(t) \le \frac{1}{(\mu_2)^{(8)}} T_{40}^0 e^{((R_1)^{(8)}+(r_{40})^{(8)})t}$ | 379 |
| $\frac{(b_{42})^{(8)}T_{40}^0}{(\mu_1)^{(8)}((R_1)^{(8)}-(b'_{42})^{(8)})}\left[e^{(R_1)^{(8)}t} - e^{-(b'_{42})^{(8)}t}\right] + T_{42}^0 e^{-(b'_{42})^{(8)}t} \le T_{42}(t) \le$ $\frac{(a_{42})^{(8)}T_{40}^0}{(\mu_2)^{(8)}((R_1)^{(8)}+(r_{40})^{(8)}+(R_2)^{(8)})}\left[e^{((R_1)^{(8)}+(r_{40})^{(8)})t} - e^{-(R_2)^{(8)}t}\right] + T_{42}^0 e^{-(R_2)^{(8)}t}$ | 380 |
| **Definition of** $(S_1)^{(8)}, (S_2)^{(8)}, (R_1)^{(8)}, (R_2)^{(8)}$:-<br>Where $(S_1)^{(8)} = (a_{40})^{(8)}(m_2)^{(8)} - (a'_{40})^{(8)}$<br>$(S_2)^{(8)} = (a_{42})^{(8)} - (p_{42})^{(8)}$<br>$(R_1)^{(8)} = (b_{40})^{(8)}(\mu_2)^{(8)} - (b'_{40})^{(8)}$<br>$(R_2)^{(8)} = (b'_{42})^{(8)} - (r_{42})^{(8)}$ | 381 |
| **Behavior of the solutions of equation 37 to 92**<br>**Theorem 2:** If we denote and define<br>**Definition of** $(\sigma_1)^{(9)}, (\sigma_2)^{(9)}, (\tau_1)^{(9)}, (\tau_2)^{(9)}$ :<br>$(\sigma_1)^{(9)}, (\sigma_2)^{(9)}, (\tau_1)^{(9)}, (\tau_2)^{(9)}$ four constants satisfying<br>$-(\sigma_2)^{(9)} \le -(a'_{44})^{(9)} + (a'_{45})^{(9)} - (a''_{44})^{(9)}(T_{45}, t) + (a''_{45})^{(9)}(T_{45}, t) \le -(\sigma_1)^{(9)}$<br>$-(\tau_2)^{(9)} \le -(b'_{44})^{(9)} + (b'_{45})^{(9)} - (b''_{44})^{(9)}((G_{47}), t) - (b''_{45})^{(9)}((G_{47}), t) \le -(\tau_1)^{(9)}$ | 382 |
| **Definition of** $(\nu_1)^{(9)}, (\nu_2)^{(9)}, (u_1)^{(9)}, (u_2)^{(9)}, \nu^{(9)}, u^{(9)}$ :<br>By $(\nu_1)^{(9)} > 0$ , $(\nu_2)^{(9)} < 0$ and respectively $(u_1)^{(9)} > 0$ , $(u_2)^{(9)} < 0$ the roots of the equations $(a_{45})^{(9)}(\nu^{(9)})^2 + (\sigma_1)^{(9)}\nu^{(9)} - (a_{44})^{(9)} = 0$<br>and $(b_{45})^{(9)}(u^{(9)})^2 + (\tau_1)^{(9)}u^{(9)} - (b_{44})^{(9)} = 0$ and | |
| **Definition of** $(\bar{\nu}_1)^{(9)},, (\bar{\nu}_2)^{(9)}, (\bar{u}_1)^{(9)}, (\bar{u}_2)^{(9)}$ : | |

By $(\bar{v}_1)^{(9)} > 0$ , $(\bar{v}_2)^{(9)} < 0$ and respectively $(\bar{u}_1)^{(9)} > 0$ , $(\bar{u}_2)^{(9)} < 0$ the roots of the equations $(a_{45})^{(9)}\left(v^{(9)}\right)^2 + (\sigma_2)^{(9)}v^{(9)} - (a_{44})^{(9)} = 0$ and $(b_{45})^{(9)}\left(u^{(9)}\right)^2 + (\tau_2)^{(9)}u^{(9)} - (b_{44})^{(9)} = 0$

**Definition of** $(m_1)^{(9)}$ , $(m_2)^{(9)}$ , $(\mu_1)^{(9)}$, $(\mu_2)^{(9)}$, $(v_0)^{(9)}$ :-

If we define $(m_1)^{(9)}$ , $(m_2)^{(9)}$ , $(\mu_1)^{(9)}$, $(\mu_2)^{(9)}$ by

$(m_2)^{(9)} = (v_0)^{(9)}, (m_1)^{(9)} = (v_1)^{(9)}$, ***if*** $(v_0)^{(9)} < (v_1)^{(9)}$

$(m_2)^{(9)} = (v_1)^{(9)}, (m_1)^{(9)} = (\bar{v}_1)^{(9)}$ , ***if*** $(v_1)^{(9)} < (v_0)^{(9)} < (\bar{v}_1)^{(9)}$,

and $\boxed{(v_0)^{(9)} = \frac{G_{44}^0}{G_{45}^0}}$

$(m_2)^{(9)} = (v_1)^{(9)}, (m_1)^{(9)} = (v_0)^{(9)}$, ***if*** $(\bar{v}_1)^{(9)} < (v_0)^{(9)}$

and analogously

$(\mu_2)^{(9)} = (u_0)^{(9)}, (\mu_1)^{(9)} = (u_1)^{(9)}$, ***if*** $(u_0)^{(9)} < (u_1)^{(9)}$

$(\mu_2)^{(9)} = (u_1)^{(9)}, (\mu_1)^{(9)} = (\bar{u}_1)^{(9)}$ , ***if*** $(u_1)^{(9)} < (u_0)^{(9)} < (\bar{u}_1)^{(9)}$,

and $\boxed{(u_0)^{(9)} = \frac{T_{44}^0}{T_{45}^0}}$

$(\mu_2)^{(9)} = (u_1)^{(9)}, (\mu_1)^{(9)} = (u_0)^{(9)}$, ***if*** $(\bar{u}_1)^{(9)} < (u_0)^{(9)}$ where $(u_1)^{(9)}, (\bar{u}_1)^{(9)}$ are defined by 59 and 69 respectively

Then the solution of 19,20,21,22,23 and 24 satisfies the inequalities

$$G_{44}^0 e^{((S_1)^{(9)} - (p_{44})^{(9)})t} \le G_{44}(t) \le G_{44}^0 e^{(S_1)^{(9)}t}$$

where $(p_i)^{(9)}$ is defined by equation 45

$$\frac{1}{(m_9)^{(9)}} G_{44}^0 e^{((S_1)^{(9)} - (p_{44})^{(9)})t} \le G_{45}(t) \le \frac{1}{(m_2)^{(9)}} G_{44}^0 e^{(S_1)^{(9)}t}$$

$$\left(\frac{(a_{46})^{(9)} G_{44}^0}{(m_1)^{(9)}((S_1)^{(9)} - (p_{44})^{(9)} - (S_2)^{(9)})}\left[e^{((S_1)^{(9)} - (p_{44})^{(9)})t} - e^{-(S_2)^{(9)}t}\right] + G_{46}^0 e^{-(S_2)^{(9)}t} \le G_{46}(t) \le \right.$$

$$\left. \frac{(a_{46})^{(9)} G_{44}^0}{(m_2)^{(9)}((S_1)^{(9)} - (a'_{46})^{(9)})}\left[e^{(S_1)^{(9)}t} - e^{-(a'_{46})^{(9)}t}\right] + G_{46}^0 e^{-(a'_{46})^{(9)}t}\right)$$

$$\boxed{T_{44}^0 e^{(R_1)^{(9)}t} \le T_{44}(t) \le T_{44}^0 e^{((R_1)^{(9)} + (r_{44})^{(9)})t}}$$

$$\frac{1}{(\mu_1)^{(9)}} T_{44}^0 e^{(R_1)^{(9)}t} \le T_{44}(t) \le \frac{1}{(\mu_2)^{(9)}} T_{44}^0 e^{((R_1)^{(9)} + (r_{44})^{(9)})t}$$

$$\frac{(b_{46})^{(9)} T_{44}^0}{(\mu_1)^{(9)}((R_1)^{(9)} - (b'_{46})^{(9)})}\left[e^{(R_1)^{(9)}t} - e^{-(b'_{46})^{(9)}t}\right] + T_{46}^0 e^{-(b'_{46})^{(9)}t} \le T_{46}(t) \le$$

$$\frac{(a_{46})^{(9)}T_{44}^{0}}{(\mu_2)^{(9)}((R_1)^{(9)}+(r_{44})^{(9)}+(R_2)^{(9)})}\left[e^{((R_1)^{(9)}+(r_{44})^{(9)})t}-e^{-(R_2)^{(9)}t}\right]+T_{46}^{0}e^{-(R_2)^{(9)}t}$$

**Definition of** $(S_1)^{(9)},(S_2)^{(9)},(R_1)^{(9)},(R_2)^{(9)}$:-

Where $(S_1)^{(9)}=(a_{44})^{(9)}(m_2)^{(9)}-(a_{44}^{'})^{(9)}$

$(S_2)^{(9)}=(a_{46})^{(9)}-(p_{46})^{(9)}$

$(R_1)^{(9)}=(b_{44})^{(9)}(\mu_2)^{(9)}-(b_{44}^{'})^{(9)}$

$(R_2)^{(9)}=(b_{46}^{'})^{(9)}-(r_{46})^{(9)}$

**Proof :** From global equations we obtain

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$$\frac{dv^{(1)}}{dt}=(a_{13})^{(1)}-\left((a_{13}^{'})^{(1)}-(a_{14}^{'})^{(1)}+(a_{13}^{''})^{(1)}(T_{14},t)\right)-(a_{14}^{''})^{(1)}(T_{14},t)v^{(1)}-(a_{14})^{(1)}v^{(1)}$$

**Definition of** $v^{(1)}$ :- $\boxed{v^{(1)}=\frac{G_{13}}{G_{14}}}$

It follows

$$-\left((a_{14})^{(1)}\left(v^{(1)}\right)^2+(\sigma_2)^{(1)}v^{(1)}-(a_{13})^{(1)}\right)\leq\frac{dv^{(1)}}{dt}\leq-\left((a_{14})^{(1)}\left(v^{(1)}\right)^2+(\sigma_1)^{(1)}v^{(1)}-(a_{13})^{(1)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(1)},(\nu_0)^{(1)}$ :-

For $0<\boxed{(\nu_0)^{(1)}=\frac{G_{13}^{0}}{G_{14}^{0}}}<(\nu_1)^{(1)}<(\bar{\nu}_1)^{(1)}$

$$\nu^{(1)}(t)\geq\frac{(\nu_1)^{(1)}+(C)^{(1)}(\nu_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)}-(\nu_0)^{(1)}\right)t\right]}}{1+(C)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)}-(\nu_0)^{(1)}\right)t\right]}}\quad,\quad\boxed{(C)^{(1)}=\frac{(\nu_1)^{(1)}-(\nu_0)^{(1)}}{(\nu_0)^{(1)}-(\nu_2)^{(1)}}}$$

it follows $(\nu_0)^{(1)}\leq\nu^{(1)}(t)\leq(\nu_1)^{(1)}$

In the same manner , we get

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$$\nu^{(1)}(t)\leq\frac{(\bar{\nu}_1)^{(1)}+(\bar{C})^{(1)}(\bar{\nu}_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}}{1+(\bar{C})^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}}\quad,\quad\boxed{(\bar{C})^{(1)}=\frac{(\bar{\nu}_1)^{(1)}-(\nu_0)^{(1)}}{(\nu_0)^{(1)}-(\bar{\nu}_2)^{(1)}}}$$

From which we deduce $(\nu_0)^{(1)}\leq\nu^{(1)}(t)\leq(\bar{\nu}_1)^{(1)}$

If $0<(\nu_1)^{(1)}<(\nu_0)^{(1)}=\frac{G_{13}^{0}}{G_{14}^{0}}<(\bar{\nu}_1)^{(1)}$ we find like in the previous case,

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$$(\nu_1)^{(1)} \le \frac{(\nu_1)^{(1)}+(C)^{(1)}(\nu_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)}-(\nu_2)^{(1)}\right)t\right]}}{1+(C)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\nu_1)^{(1)}-(\nu_2)^{(1)}\right)t\right]}} \le \nu^{(1)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(1)}+(\bar{C})^{(1)}(\bar{\nu}_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}}{1+(\bar{C})^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}} \le (\bar{\nu}_1)^{(1)}$$

If $0 < (\nu_1)^{(1)} \le (\bar{\nu}_1)^{(1)} \le \boxed{(\nu_0)^{(1)} = \frac{G_{13}^0}{G_{14}^0}}$ , we obtain 386

$$(\nu_1)^{(1)} \le \nu^{(1)}(t) \le \frac{(\bar{\nu}_1)^{(1)}+(\bar{C})^{(1)}(\bar{\nu}_2)^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}}{1+(\bar{C})^{(1)}e^{\left[-(a_{14})^{(1)}\left((\bar{\nu}_1)^{(1)}-(\bar{\nu}_2)^{(1)}\right)t\right]}} \le (\nu_0)^{(1)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(1)}(t)$ :-

$$(m_2)^{(1)} \le \nu^{(1)}(t) \le (m_1)^{(1)}, \quad \boxed{\nu^{(1)}(t) = \frac{G_{13}(t)}{G_{14}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(1)}(t)$ :-

$$(\mu_2)^{(1)} \le u^{(1)}(t) \le (\mu_1)^{(1)}, \quad \boxed{u^{(1)}(t) = \frac{T_{13}(t)}{T_{14}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{13}^{''})^{(1)} = (a_{14}^{''})^{(1)}, then\ (\sigma_1)^{(1)} = (\sigma_2)^{(1)}$ and in this case $(\nu_1)^{(1)} = (\bar{\nu}_1)^{(1)}$ if in addition $(\nu_0)^{(1)} = (\nu_1)^{(1)}$ then $\nu^{(1)}(t) = (\nu_0)^{(1)}$ and as a consequence $G_{13}(t) = (\nu_0)^{(1)}G_{14}(t)$ this also defines $(\nu_0)^{(1)}$ for the special case

Analogously if $(b_{13}^{''})^{(1)} = (b_{14}^{''})^{(1)}, then\ (\tau_1)^{(1)} = (\tau_2)^{(1)}$ and then

$(u_1)^{(1)} = (\bar{u}_1)^{(1)}$if in addition $(u_0)^{(1)} = (u_1)^{(1)}$ then $T_{13}(t) = (u_0)^{(1)}T_{14}(t)$ This is an important consequence of the relation between $(\nu_1)^{(1)}$ and $(\bar{\nu}_1)^{(1)}$, and definition of $(u_0)^{(1)}$.

**Proof** : From global equations we obtain 387

$$\frac{d\nu^{(2)}}{dt} = (a_{16})^{(2)} - \left((a_{16}')^{(2)} - (a_{17}')^{(2)} + (a_{16}'')^{(2)}(T_{17}, t)\right) - (a_{17}'')^{(2)}(T_{17}, t)\nu^{(2)} - (a_{17})^{(2)}\nu^{(2)}$$

**Definition of** $\nu^{(2)}$ :- $\boxed{\nu^{(2)} = \frac{G_{16}}{G_{17}}}$ 388

It follows 389

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| $-\left((a_{17})^{(2)}\left(v^{(2)}\right)^2+(\sigma_2)^{(2)}v^{(2)}-(a_{16})^{(2)}\right)\le\frac{\mathrm{d}v^{(2)}}{\mathrm{dt}}\le-\left((a_{17})^{(2)}\left(v^{(2)}\right)^2+(\sigma_1)^{(2)}v^{(2)}-(a_{16})^{(2)}\right)$ | |
| From which one obtains<br><br>**Definition of** $(\bar{v}_1)^{(2)},(v_0)^{(2)}$ :-<br><br>For $0<(v_0)^{(2)}=\frac{G_{16}^0}{G_{17}^0}<(v_1)^{(2)}<(\bar{v}_1)^{(2)}$<br><br>$v^{(2)}(t)\ge\frac{(v_1)^{(2)}+(C)^{(2)}(v_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((v_1)^{(2)}-(v_0)^{(2)}\right)t\right]}}{1+(C)^{(2)}e^{\left[-(a_{17})^{(2)}\left((v_1)^{(2)}-(v_0)^{(2)}\right)t\right]}}$ , $\boxed{(C)^{(2)}=\frac{(v_1)^{(2)}-(v_0)^{(2)}}{(v_0)^{(2)}-(v_2)^{(2)}}}$<br><br>it follows $(v_0)^{(2)}\le v^{(2)}(t)\le(v_1)^{(2)}$ | 390 |
| In the same manner , we get<br><br>$v^{(2)}(t)\le\frac{(\bar{v}_1)^{(2)}+(\bar{C})^{(2)}(\bar{v}_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}{1+(\bar{C})^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}$ , $\boxed{(\bar{C})^{(2)}=\frac{(\bar{v}_1)^{(2)}-(v_0)^{(2)}}{(v_0)^{(2)}-(\bar{v}_2)^{(2)}}}$ | 391 |
| From which we deduce $(v_0)^{(2)}\le v^{(2)}(t)\le(\bar{v}_1)^{(2)}$ | 392 |
| If $0<(v_1)^{(2)}<(v_0)^{(2)}=\frac{G_{16}^0}{G_{17}^0}<(\bar{v}_1)^{(2)}$ we find like in the previous case,<br><br>$(v_1)^{(2)}\le\frac{(v_1)^{(2)}+(C)^{(2)}(v_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((v_1)^{(2)}-(v_2)^{(2)}\right)t\right]}}{1+(C)^{(2)}e^{\left[-(a_{17})^{(2)}\left((v_1)^{(2)}-(v_2)^{(2)}\right)t\right]}}\le v^{(2)}(t)\le$<br><br>$\frac{(\bar{v}_1)^{(2)}+(\bar{C})^{(2)}(\bar{v}_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}{1+(\bar{C})^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}\le(\bar{v}_1)^{(2)}$ | 393 |
| If $0<(v_1)^{(2)}\le(\bar{v}_1)^{(2)}\le(v_0)^{(2)}=\frac{G_{16}^0}{G_{17}^0}$ , we obtain<br><br>$(v_1)^{(2)}\le v^{(2)}(t)\le\frac{(\bar{v}_1)^{(2)}+(\bar{C})^{(2)}(\bar{v}_2)^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}{1+(\bar{C})^{(2)}e^{\left[-(a_{17})^{(2)}\left((\bar{v}_1)^{(2)}-(\bar{v}_2)^{(2)}\right)t\right]}}\le(v_0)^{(2)}$<br><br>And so with the notation of the first part of condition (c) , we have | 394 |
| **Definition of** $v^{(2)}(t)$ :-<br><br>$(m_2)^{(2)}\le v^{(2)}(t)\le(m_1)^{(2)}$, $\boxed{v^{(2)}(t)=\frac{G_{16}(t)}{G_{17}(t)}}$ | 395 |
| In a completely analogous way, we obtain<br><br>**Definition of** $u^{(2)}(t)$ :- | 396 |

| | |
|---|---|
| $(\mu_2)^{(2)} \leq u^{(2)}(t) \leq (\mu_1)^{(2)}$, $\boxed{u^{(2)}(t) = \frac{T_{16}(t)}{T_{17}(t)}}$ | |
| Now, using this result and replacing it in global equations we get easily the result stated in the theorem. | |
| **Particular case :**<br><br>If $(a_{16}^{''})^{(2)} = (a_{17}^{''})^{(2)}, then\ (\sigma_1)^{(2)} = (\sigma_2)^{(2)}$ and in this case $(\nu_1)^{(2)} = (\bar{\nu}_1)^{(2)}$ if in addition $(\nu_0)^{(2)} = (\nu_1)^{(2)}$ then $\nu^{(2)}(t) = (\nu_0)^{(2)}$ and as a consequence $G_{16}(t) = (\nu_0)^{(2)} G_{17}(t)$<br><br>Analogously if $(b_{16}^{''})^{(2)} = (b_{17}^{''})^{(2)}, then\ (\tau_1)^{(2)} = (\tau_2)^{(2)}$ and then<br><br>$(u_1)^{(2)} = (\bar{u}_1)^{(2)}$if in addition $(u_0)^{(2)} = (u_1)^{(2)}$ then $T_{16}(t) = (u_0)^{(2)} T_{17}(t)$ This is an important consequence of the relation between $(\nu_1)^{(2)}$ and $(\bar{\nu}_1)^{(2)}$ | 397 |
| **Proof** : From global equations we obtain<br><br>$\frac{d\nu^{(3)}}{dt} = (a_{20})^{(3)} - \left((a_{20}^{'})^{(3)} - (a_{21}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21}, t)\right) - (a_{21}^{''})^{(3)}(T_{21}, t)\nu^{(3)} - (a_{21})^{(3)}\nu^{(3)}$ | 398 |
| **Definition of** $\nu^{(3)}$ :- $\boxed{\nu^{(3)} = \frac{G_{20}}{G_{21}}}$<br><br>It follows<br><br>$-\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_2)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right) \leq \frac{d\nu^{(3)}}{dt} \leq -\left((a_{21})^{(3)}\left(\nu^{(3)}\right)^2 + (\sigma_1)^{(3)}\nu^{(3)} - (a_{20})^{(3)}\right)$ | 399 |
| From which one obtains<br><br>For $0 < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\nu_1)^{(3)} < (\bar{\nu}_1)^{(3)}$<br><br>$\nu^{(3)}(t) \geq \frac{(\nu_1)^{(3)} + (C)^{(3)}(\nu_2)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)} - (\nu_0)^{(3)}\right)t\right]}}{1 + (C)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)} - (\nu_0)^{(3)}\right)t\right]}}$ , $\boxed{(C)^{(3)} = \frac{(\nu_1)^{(3)} - (\nu_0)^{(3)}}{(\nu_0)^{(3)} - (\nu_2)^{(3)}}}$<br><br>it follows $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\nu_1)^{(3)}$ | 400 |
| In the same manner , we get<br><br>$\nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)} + (C)^{(3)}(\bar{\nu}_2)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)} - (\bar{\nu}_2)^{(3)}\right)t\right]}}{1 + (C)^{(3)} e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)} - (\bar{\nu}_2)^{(3)}\right)t\right]}}$ , $\boxed{(\bar{C})^{(3)} = \frac{(\bar{\nu}_1)^{(3)} - (\nu_0)^{(3)}}{(\nu_0)^{(3)} - (\bar{\nu}_2)^{(3)}}}$<br><br>**Definition of** $(\bar{\nu}_1)^{(3)}$ :-<br><br>From which we deduce $(\nu_0)^{(3)} \leq \nu^{(3)}(t) \leq (\bar{\nu}_1)^{(3)}$ | 401 |

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If $0 < (\nu_1)^{(3)} < (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0} < (\bar{\nu}_1)^{(3)}$ we find like in the previous case,

$$(\nu_1)^{(3)} \leq \frac{(\nu_1)^{(3)}+(C)^{(3)}(\nu_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}}{1+(C)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\nu_1)^{(3)}-(\nu_2)^{(3)}\right)t\right]}} \leq \nu^{(3)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(C)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\bar{\nu}_1)^{(3)}$$

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If $0 < (\nu_1)^{(3)} \leq (\bar{\nu}_1)^{(3)} \leq (\nu_0)^{(3)} = \frac{G_{20}^0}{G_{21}^0}$ , we obtain

$$(\nu_1)^{(3)} \leq \nu^{(3)}(t) \leq \frac{(\bar{\nu}_1)^{(3)}+(\bar{C})^{(3)}(\bar{\nu}_2)^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}}{1+(\bar{C})^{(3)}e^{\left[-(a_{21})^{(3)}\left((\bar{\nu}_1)^{(3)}-(\bar{\nu}_2)^{(3)}\right)t\right]}} \leq (\nu_0)^{(3)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(3)}(t)$ :-

$$(m_2)^{(3)} \leq \nu^{(3)}(t) \leq (m_1)^{(3)}, \quad \boxed{\nu^{(3)}(t) = \frac{G_{20}(t)}{G_{21}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(3)}(t)$ :-

$$(\mu_2)^{(3)} \leq u^{(3)}(t) \leq (\mu_1)^{(3)}, \quad \boxed{u^{(3)}(t) = \frac{T_{20}(t)}{T_{21}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

Particular case :

If $(a_{20}'')^{(3)} = (a_{21}'')^{(3)}, then\ (\sigma_1)^{(3)} = (\sigma_2)^{(3)}$ and in this case $(\nu_1)^{(3)} = (\bar{\nu}_1)^{(3)}$ if in addition $(\nu_0)^{(3)} = (\nu_1)^{(3)}$ then $\nu^{(3)}(t) = (\nu_0)^{(3)}$ and as a consequence $G_{20}(t) = (\nu_0)^{(3)}G_{21}(t)$

Analogously if $(b_{20}'')^{(3)} = (b_{21}'')^{(3)}, then\ (\tau_1)^{(3)} = (\tau_2)^{(3)}$ and then

$(u_1)^{(3)} = (\bar{u}_1)^{(3)}$if in addition $(u_0)^{(3)} = (u_1)^{(3)}$ then $T_{20}(t) = (u_0)^{(3)}T_{21}(t)$ This is an important consequence of the relation between $(\nu_1)^{(3)}$ and $(\bar{\nu}_1)^{(3)}$

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(4)}}{dt} = (a_{24})^{(4)} - \left((a_{24}')^{(4)} - (a_{25}')^{(4)} + (a_{24}'')^{(4)}(T_{25},t)\right) - (a_{25}'')^{(4)}(T_{25},t)\nu^{(4)} - (a_{25})^{(4)}\nu^{(4)}$$

**Definition of** $\nu^{(4)}$ :- $\boxed{\nu^{(4)} = \frac{G_{24}}{G_{25}}}$

It follows

$$-\left((a_{25})^{(4)}\left(v^{(4)}\right)^2+(\sigma_2)^{(4)}v^{(4)}-(a_{24})^{(4)}\right)\le\frac{dv^{(4)}}{dt}\le-\left((a_{25})^{(4)}\left(v^{(4)}\right)^2+(\sigma_4)^{(4)}v^{(4)}-(a_{24})^{(4)}\right)$$

From which one obtains

**Definition of** $(\bar{v}_1)^{(4)},(v_0)^{(4)}$ :-

For $0<\boxed{(v_0)^{(4)}=\frac{G_{24}^0}{G_{25}^0}}<(v_1)^{(4)}<(\bar{v}_1)^{(4)}$

$$v^{(4)}(t)\ge\frac{(v_1)^{(4)}+(C)^{(4)}(v_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((v_1)^{(4)}-(v_0)^{(4)}\right)t\right]}}{4+(C)^{(4)}e^{\left[-(a_{25})^{(4)}\left((v_1)^{(4)}-(v_0)^{(4)}\right)t\right]}}\quad,\quad\boxed{(C)^{(4)}=\frac{(v_1)^{(4)}-(v_0)^{(4)}}{(v_0)^{(4)}-(v_2)^{(4)}}}$$

it follows $(v_0)^{(4)}\le v^{(4)}(t)\le(v_1)^{(4)}$

In the same manner , we get

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$$v^{(4)}(t)\le\frac{(\bar{v}_1)^{(4)}+(\bar{C})^{(4)}(\bar{v}_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}{4+(\bar{C})^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}\quad,\quad\boxed{(\bar{C})^{(4)}=\frac{(\bar{v}_1)^{(4)}-(v_0)^{(4)}}{(v_0)^{(4)}-(\bar{v}_2)^{(4)}}}$$

From which we deduce $(v_0)^{(4)}\le v^{(4)}(t)\le(\bar{v}_1)^{(4)}$

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If $0<(v_1)^{(4)}<(v_0)^{(4)}=\frac{G_{24}^0}{G_{25}^0}<(\bar{v}_1)^{(4)}$ we find like in the previous case,

$$(v_1)^{(4)}\le\frac{(v_1)^{(4)}+(C)^{(4)}(v_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((v_1)^{(4)}-(v_2)^{(4)}\right)t\right]}}{1+(C)^{(4)}e^{\left[-(a_{25})^{(4)}\left((v_1)^{(4)}-(v_2)^{(4)}\right)t\right]}}\le v^{(4)}(t)\le$$

$$\frac{(\bar{v}_1)^{(4)}+(\bar{C})^{(4)}(\bar{v}_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}{1+(\bar{C})^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}\le(\bar{v}_1)^{(4)}$$

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If $0<(v_1)^{(4)}\le(\bar{v}_1)^{(4)}\le\boxed{(v_0)^{(4)}=\frac{G_{24}^0}{G_{25}^0}}$ , we obtain

$$(v_1)^{(4)}\le v^{(4)}(t)\le\frac{(\bar{v}_1)^{(4)}+(\bar{C})^{(4)}(\bar{v}_2)^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}{1+(\bar{C})^{(4)}e^{\left[-(a_{25})^{(4)}\left((\bar{v}_1)^{(4)}-(\bar{v}_2)^{(4)}\right)t\right]}}\le(v_0)^{(4)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $v^{(4)}(t)$ :-

$$(m_2)^{(4)}\le v^{(4)}(t)\le(m_1)^{(4)},\quad\boxed{v^{(4)}(t)=\frac{G_{24}(t)}{G_{25}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(4)}(t)$ :-

$$(\mu_2)^{(4)}\le u^{(4)}(t)\le(\mu_1)^{(4)},\quad\boxed{u^{(4)}(t)=\frac{T_{24}(t)}{T_{25}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{24}^{''})^{(4)} = (a_{25}^{''})^{(4)}, then\ (\sigma_1)^{(4)} = (\sigma_2)^{(4)}$ and in this case $(\nu_1)^{(4)} = (\bar{\nu}_1)^{(4)}$ if in addition $(\nu_0)^{(4)} = (\nu_1)^{(4)}$ then $\nu^{(4)}(t) = (\nu_0)^{(4)}$ and as a consequence $G_{24}(t) = (\nu_0)^{(4)} G_{25}(t)$ **this also defines** $(\nu_0)^{(4)}$ **for the special case** .

Analogously if $(b_{24}^{''})^{(4)} = (b_{25}^{''})^{(4)}, then\ (\tau_1)^{(4)} = (\tau_2)^{(4)}$ and then $(u_1)^{(4)} = (\bar{u}_4)^{(4)}$ if in addition $(u_0)^{(4)} = (u_1)^{(4)}$ then $T_{24}(t) = (u_0)^{(4)} T_{25}(t)$ This is an important consequence of the relation between $(\nu_1)^{(4)}$ and $(\bar{\nu}_1)^{(4)}$, **and definition of** $(u_0)^{(4)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(5)}}{dt} = (a_{28})^{(5)} - \left((a_{28}^{'})^{(5)} - (a_{29}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}, t)\right) - (a_{29}^{''})^{(5)}(T_{29}, t)\nu^{(5)} - (a_{29})^{(5)}\nu^{(5)}$$

**Definition of** $\nu^{(5)}$ :- $\boxed{\nu^{(5)} = \frac{G_{28}}{G_{29}}}$

It follows

$$-\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_2)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right) \leq \frac{d\nu^{(5)}}{dt} \leq -\left((a_{29})^{(5)}\left(\nu^{(5)}\right)^2 + (\sigma_1)^{(5)}\nu^{(5)} - (a_{28})^{(5)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(5)}, (\nu_0)^{(5)}$ :-

For $0 < \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}} < (\nu_1)^{(5)} < (\bar{\nu}_1)^{(5)}$

$$\nu^{(5)}(t) \geq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}}{5 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_0)^{(5)}\right)t\right]}} \quad , \quad \boxed{(C)^{(5)} = \frac{(\nu_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\nu_2)^{(5)}}}$$

it follows $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\nu_1)^{(5)}$

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In the same manner , we get

$$\nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)} + (\bar{C})^{(5)}(\bar{\nu}_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}}{5 + (\bar{C})^{(5)} e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)} - (\bar{\nu}_2)^{(5)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(5)} = \frac{(\bar{\nu}_1)^{(5)} - (\nu_0)^{(5)}}{(\nu_0)^{(5)} - (\bar{\nu}_2)^{(5)}}}$$

From which we deduce $(\nu_0)^{(5)} \leq \nu^{(5)}(t) \leq (\bar{\nu}_5)^{(5)}$

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If $0 < (\nu_1)^{(5)} < (\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0} < (\bar{\nu}_1)^{(5)}$ we find like in the previous case,

$$(\nu_1)^{(5)} \leq \frac{(\nu_1)^{(5)} + (C)^{(5)}(\nu_2)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}}{1 + (C)^{(5)} e^{\left[-(a_{29})^{(5)}\left((\nu_1)^{(5)} - (\nu_2)^{(5)}\right)t\right]}} \leq \nu^{(5)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(5)}+(\bar{C})^{(5)}(\bar{\nu}_2)^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}}{1+(\bar{C})^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\bar{\nu}_1)^{(5)}$$

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If $0 < (\nu_1)^{(5)} \leq (\bar{\nu}_1)^{(5)} \leq \boxed{(\nu_0)^{(5)} = \frac{G_{28}^0}{G_{29}^0}}$ , we obtain

$$(\nu_1)^{(5)} \leq \nu^{(5)}(t) \leq \frac{(\bar{\nu}_1)^{(5)}+(\bar{C})^{(5)}(\bar{\nu}_2)^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}}{1+(\bar{C})^{(5)}e^{\left[-(a_{29})^{(5)}\left((\bar{\nu}_1)^{(5)}-(\bar{\nu}_2)^{(5)}\right)t\right]}} \leq (\nu_0)^{(5)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(5)}(t)$ :-

$$(m_2)^{(5)} \leq \nu^{(5)}(t) \leq (m_1)^{(5)}, \quad \boxed{\nu^{(5)}(t) = \frac{G_{28}(t)}{G_{29}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(5)}(t)$ :-

$$(\mu_2)^{(5)} \leq u^{(5)}(t) \leq (\mu_1)^{(5)}, \quad \boxed{u^{(5)}(t) = \frac{T_{28}(t)}{T_{29}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{28}^{''})^{(5)} = (a_{29}^{''})^{(5)}, then\ (\sigma_1)^{(5)} = (\sigma_2)^{(5)}$ and in this case $(\nu_1)^{(5)} = (\bar{\nu}_1)^{(5)}$ if in addition $(\nu_0)^{(5)} = (\nu_5)^{(5)}$ then $\nu^{(5)}(t) = (\nu_0)^{(5)}$ and as a consequence $G_{28}(t) = (\nu_0)^{(5)}G_{29}(t)$ **this also defines** $(\nu_0)^{(5)}$ **for the special case** .

Analogously if $(b_{28}^{''})^{(5)} = (b_{29}^{''})^{(5)}, then\ (\tau_1)^{(5)} = (\tau_2)^{(5)}$ and then $(u_1)^{(5)} = (\bar{u}_1)^{(5)}$ if in addition $(u_0)^{(5)} = (u_1)^{(5)}$ then $T_{28}(t) = (u_0)^{(5)}T_{29}(t)$ This is an important consequence of the relation between $(\nu_1)^{(5)}$ and $(\bar{\nu}_1)^{(5)}$, **and definition of** $(u_0)^{(5)}$.

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**Proof** : From global equations we obtain

$$\frac{d\nu^{(6)}}{dt} = (a_{32})^{(6)} - \left((a_{32}^{'})^{(6)} - (a_{33}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33}, t)\right) - (a_{33}^{''})^{(6)}(T_{33}, t)\nu^{(6)} - (a_{33})^{(6)}\nu^{(6)}$$

**Definition of** $\nu^{(6)}$ :- $\boxed{\nu^{(6)} = \frac{G_{32}}{G_{33}}}$

It follows

$$-\left((a_{33})^{(6)}\left(\nu^{(6)}\right)^2 + (\sigma_2)^{(6)}\nu^{(6)} - (a_{32})^{(6)}\right) \leq \frac{d\nu^{(6)}}{dt} \leq -\left((a_{33})^{(6)}\left(\nu^{(6)}\right)^2 + (\sigma_1)^{(6)}\nu^{(6)} - (a_{32})^{(6)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(6)}, (\nu_0)^{(6)}$ :-

For $0 < \boxed{(\nu_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0}} < (\nu_1)^{(6)} < (\bar{\nu}_1)^{(6)}$

$$\nu^{(6)}(t) \geq \frac{(\nu_1)^{(6)}+(C)^{(6)}(\nu_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_0)^{(6)}\right)t\right]}}{1+(C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_0)^{(6)}\right)t\right]}} \quad , \quad \boxed{(C)^{(6)}=\frac{(\nu_1)^{(6)}-(\nu_0)^{(6)}}{(\nu_0)^{(6)}-(\nu_2)^{(6)}}}$$

it follows $(\nu_0)^{(6)} \leq \nu^{(6)}(t) \leq (\nu_1)^{(6)}$

In the same manner , we get

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$$\nu^{(6)}(t) \leq \frac{(\bar{\nu}_1)^{(6)}+(\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}}{1+(\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(6)}=\frac{(\bar{\nu}_1)^{(6)}-(\nu_0)^{(6)}}{(\nu_0)^{(6)}-(\bar{\nu}_2)^{(6)}}}$$

From which we deduce $(\nu_0)^{(6)} \leq \nu^{(6)}(t) \leq (\bar{\nu}_1)^{(6)}$

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If $0 < (\nu_1)^{(6)} < (\nu_0)^{(6)} = \frac{G_{32}^0}{G_{33}^0} < (\bar{\nu}_1)^{(6)}$ we find like in the previous case,

$$(\nu_1)^{(6)} \leq \frac{(\nu_1)^{(6)}+(C)^{(6)}(\nu_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_2)^{(6)}\right)t\right]}}{1+(C)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\nu_1)^{(6)}-(\nu_2)^{(6)}\right)t\right]}} \leq \nu^{(6)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(6)}+(\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}}{1+(\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}} \leq (\bar{\nu}_1)^{(6)}$$

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If $0 < (\nu_1)^{(6)} \leq (\bar{\nu}_1)^{(6)} \leq \boxed{(\nu_0)^{(6)}=\frac{G_{32}^0}{G_{33}^0}}$ , we obtain

$$(\nu_1)^{(6)} \leq \nu^{(6)}(t) \leq \frac{(\bar{\nu}_1)^{(6)}+(\bar{C})^{(6)}(\bar{\nu}_2)^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}}{1+(\bar{C})^{(6)}e^{\left[-(a_{33})^{(6)}\left((\bar{\nu}_1)^{(6)}-(\bar{\nu}_2)^{(6)}\right)t\right]}} \leq (\nu_0)^{(6)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(6)}(t)$ :-

$$(m_2)^{(6)} \leq \nu^{(6)}(t) \leq (m_1)^{(6)}, \quad \boxed{\nu^{(6)}(t)=\frac{G_{32}(t)}{G_{33}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(6)}(t)$ :-

$$(\mu_2)^{(6)} \leq u^{(6)}(t) \leq (\mu_1)^{(6)}, \quad \boxed{u^{(6)}(t)=\frac{T_{32}(t)}{T_{33}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{32}^{''})^{(6)}=(a_{33}^{''})^{(6)}, then\ (\sigma_1)^{(6)}=(\sigma_2)^{(6)}$ and in this case $(\nu_1)^{(6)}=(\bar{\nu}_1)^{(6)}$ if in addition $(\nu_0)^{(6)}=(\nu_1)^{(6)}$ then $\nu^{(6)}(t)=(\nu_0)^{(6)}$ and as a consequence $G_{32}(t)=(\nu_0)^{(6)}G_{33}(t)$ **this also defines** $(\nu_0)^{(6)}$ **for the special case** .

Analogously if $(b_{32}^{''})^{(6)}=(b_{33}^{''})^{(6)}, then\ (\tau_1)^{(6)}=(\tau_2)^{(6)}$ and then $(u_1)^{(6)}=(\bar{u}_1)^{(6)}$ if in addition $(u_0)^{(6)}=(u_1)^{(6)}$ then $T_{32}(t)=(u_0)^{(6)}T_{33}(t)$ This is an important consequence of the relation between $(\nu_1)^{(6)}$ and $(\bar{\nu}_1)^{(6)}$, **and definition of** $(u_0)^{(6)}$.

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**Proof** : From global equations we obtain

$$\frac{dv^{(7)}}{dt} = (a_{36})^{(7)} - \left((a'_{36})^{(7)} - (a'_{37})^{(7)} + (a''_{36})^{(7)}(T_{37},t)\right) - (a''_{37})^{(7)}(T_{37},t)v^{(7)} - (a_{37})^{(7)}v^{(7)}$$

**Definition of** $v^{(7)}$ :- $\boxed{v^{(7)} = \frac{G_{36}}{G_{37}}}$

It follows

$$-\left((a_{37})^{(7)}\left(v^{(7)}\right)^2 + (\sigma_2)^{(7)}v^{(7)} - (a_{36})^{(7)}\right) \le \frac{dv^{(7)}}{dt} \le -\left((a_{37})^{(7)}\left(v^{(7)}\right)^2 + (\sigma_1)^{(7)}v^{(7)} - (a_{36})^{(7)}\right)$$

From which one obtains

**Definition of** $(\bar{v}_1)^{(7)}, (v_0)^{(7)}$ :-

For $0 < \boxed{(v_0)^{(7)} = \frac{G^0_{36}}{G^0_{37}}} < (v_1)^{(7)} < (\bar{v}_1)^{(7)}$

$$v^{(7)}(t) \ge \frac{(v_1)^{(7)} + (C)^{(7)}(v_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((v_1)^{(7)} - (v_0)^{(7)}\right)t\right]}}{1 + (C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((v_1)^{(7)} - (v_0)^{(7)}\right)t\right]}} \quad , \quad \boxed{(C)^{(7)} = \frac{(v_1)^{(7)} - (v_0)^{(7)}}{(v_0)^{(7)} - (v_2)^{(7)}}}$$

it follows $(v_0)^{(7)} \le v^{(7)}(t) \le (v_1)^{(7)}$

In the same manner , we get

$$v^{(7)}(t) \le \frac{(\bar{v}_1)^{(7)} + (\bar{C})^{(7)}(\bar{v}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}}{1 + (\bar{C})^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(7)} = \frac{(\bar{v}_1)^{(7)} - (v_0)^{(7)}}{(v_0)^{(7)} - (\bar{v}_2)^{(7)}}}$$

From which we deduce $(v_0)^{(7)} \le v^{(7)}(t) \le (\bar{v}_1)^{(7)}$

If $0 < (v_1)^{(7)} < (v_0)^{(7)} = \frac{G^0_{36}}{G^0_{37}} < (\bar{v}_1)^{(7)}$ we find like in the previous case,

$$(v_1)^{(7)} \le \frac{(v_1)^{(7)} + (C)^{(7)}(v_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((v_1)^{(7)} - (v_2)^{(7)}\right)t\right]}}{1 + (C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((v_1)^{(7)} - (v_2)^{(7)}\right)t\right]}} \le v^{(7)}(t) \le$$

$$\frac{(\bar{v}_1)^{(7)} + (\bar{C})^{(7)}(\bar{v}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}}{1 + (C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}} \le (\bar{v}_1)^{(7)}$$

If $0 < (v_1)^{(7)} \le (\bar{v}_1)^{(7)} \le \boxed{(v_0)^{(7)} = \frac{G^0_{36}}{G^0_{37}}}$ , we obtain

$$(v_1)^{(7)} \le v^{(7)}(t) \le \frac{(\bar{v}_1)^{(7)} + (\bar{C})^{(7)}(\bar{v}_2)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}}{1 + (C)^{(7)}e^{\left[-(a_{37})^{(7)}\left((\bar{v}_1)^{(7)} - (\bar{v}_2)^{(7)}\right)t\right]}} \le (v_0)^{(7)}$$

And so with the notation of the first part of condition (c) , we have

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**Definition of** $v^{(7)}(t)$ :-

$(m_2)^{(7)} \leq v^{(7)}(t) \leq (m_1)^{(7)}$, $\boxed{v^{(7)}(t) = \frac{G_{36}(t)}{G_{37}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(7)}(t)$ :-

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$(\mu_2)^{(7)} \leq u^{(7)}(t) \leq (\mu_1)^{(7)}$, $\boxed{u^{(7)}(t) = \frac{T_{36}(t)}{T_{37}(t)}}$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{36}^{''})^{(7)} = (a_{37}^{''})^{(7)}, then\ (\sigma_1)^{(7)} = (\sigma_2)^{(7)}$ and in this case $(\nu_1)^{(7)} = (\bar{\nu}_1)^{(7)}$ if in addition $(\nu_0)^{(7)} = (\nu_1)^{(7)}$ then $v^{(7)}(t) = (\nu_0)^{(7)}$ and as a consequence $G_{36}(t) = (\nu_0)^{(7)} G_{37}(t)$ **this also defines** $(\nu_0)^{(7)}$ **for the special case** .

Analogously if $(b_{36}^{''})^{(7)} = (b_{37}^{''})^{(7)}, then\ (\tau_1)^{(7)} = (\tau_2)^{(7)}$ and then $(u_1)^{(7)} = (\bar{u}_1)^{(7)}$if in addition $(u_0)^{(7)} = (u_1)^{(7)}$ then $T_{36}(t) = (u_0)^{(7)} T_{37}(t)$ This is an important consequence of the relation between $(\nu_1)^{(7)}$ and $(\bar{\nu}_1)^{(7)}$, **and definition of** $(u_0)^{(7)}$.

**Proof** : From global equations we obtain

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$$\frac{dv^{(8)}}{dt} = (a_{40})^{(8)} - \left((a_{40}^{'})^{(8)} - (a_{41}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41}, t)\right) - (a_{41}^{''})^{(8)}(T_{41}, t) v^{(8)} - (a_{41})^{(8)} v^{(8)}$$

**Definition of** $v^{(8)}$ :- $\boxed{v^{(8)} = \frac{G_{40}}{G_{41}}}$

It follows

$$-\left((a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_2)^{(8)} v^{(8)} - (a_{40})^{(8)}\right) \leq \frac{dv^{(8)}}{dt} \leq -\left((a_{41})^{(8)}\left(v^{(8)}\right)^2 + (\sigma_1)^{(8)} v^{(8)} - (a_{40})^{(8)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(8)}, (\nu_0)^{(8)}$ :-

For $0 < \boxed{(\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}} < (\nu_1)^{(8)} < (\bar{\nu}_1)^{(8)}$

$$v^{(8)}(t) \geq \frac{(\nu_1)^{(8)} + (C)^{(8)}(\nu_2)^{(8)} e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}}{1 + (C)^{(8)} e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)} - (\nu_0)^{(8)}\right)t\right]}} \quad , \quad \boxed{(C)^{(8)} = \frac{(\nu_1)^{(8)} - (\nu_0)^{(8)}}{(\nu_0)^{(8)} - (\nu_2)^{(8)}}}$$

it follows $(\nu_0)^{(8)} \leq \nu^{(8)}(t) \leq (\nu_1)^{(8)}$

In the same manner , we get

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$$\nu^{(8)}(t) \leq \frac{(\bar{\nu}_1)^{(8)}+(\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}}{1+(\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(8)} = \frac{(\bar{\nu}_1)^{(8)}-(\nu_0)^{(8)}}{(\nu_0)^{(8)}-(\bar{\nu}_2)^{(8)}}}$$

From which we deduce $(\nu_0)^{(8)} \leq \nu^{(8)}(t) \leq (\bar{\nu}_8)^{(8)}$

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If $0 < (\nu_1)^{(8)} < (\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0} < (\bar{\nu}_1)^{(8)}$ we find like in the previous case,

$$(\nu_1)^{(8)} \leq \frac{(\nu_1)^{(8)}+(C)^{(8)}(\nu_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)}-(\nu_2)^{(8)}\right)t\right]}}{1+(C)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\nu_1)^{(8)}-(\nu_2)^{(8)}\right)t\right]}} \leq \nu^{(8)}(t) \leq$$

$$\frac{(\bar{\nu}_1)^{(8)}+(\bar{C})^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}}{1+(\bar{C})^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}} \leq (\bar{\nu}_1)^{(8)}$$

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If $0 < (\nu_1)^{(8)} \leq (\bar{\nu}_1)^{(8)} \leq \boxed{(\nu_0)^{(8)} = \frac{G_{40}^0}{G_{41}^0}}$ , we obtain

$$(\nu_1)^{(8)} \leq \nu^{(8)}(t) \leq \frac{(\bar{\nu}_1)^{(8)}+(C)^{(8)}(\bar{\nu}_2)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}}{1+(C)^{(8)}e^{\left[-(a_{41})^{(8)}\left((\bar{\nu}_1)^{(8)}-(\bar{\nu}_2)^{(8)}\right)t\right]}} \leq (\nu_0)^{(8)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(8)}(t)$ :-

$$(m_2)^{(8)} \leq \nu^{(8)}(t) \leq (m_1)^{(8)}, \quad \boxed{\nu^{(8)}(t) = \frac{G_{40}(t)}{G_{41}(t)}}$$

In a completely analogous way, we obtain

**Definition of** $u^{(8)}(t)$ :-

$$(\mu_2)^{(8)} \leq u^{(8)}(t) \leq (\mu_1)^{(8)}, \quad \boxed{u^{(8)}(t) = \frac{T_{40}(t)}{T_{41}(t)}}$$

Now, using this result and replacing it in global equations we get easily the result stated in the theorem.

**Particular case :**

If $(a_{40}^{''})^{(8)} = (a_{41}^{''})^{(8)}, then\ (\sigma_1)^{(8)} = (\sigma_2)^{(8)}$ and in this case $(\nu_1)^{(8)} = (\bar{\nu}_1)^{(8)}$ if in addition $(\nu_0)^{(8)} = (\nu_1)^{(8)}$ then $\nu^{(8)}(t) = (\nu_0)^{(8)}$ and as a consequence $G_{40}(t) = (\nu_0)^{(8)}G_{41}(t)$ **this also defines** $(\nu_0)^{(8)}$ **for the special case** .

Analogously if $(b_{40}^{''})^{(8)} = (b_{41}^{''})^{(8)}, then\ (\tau_1)^{(8)} = (\tau_2)^{(8)}$ and then $(u_1)^{(8)} = (\bar{u}_1)^{(8)}$ if in addition $(u_0)^{(8)} = (u_1)^{(8)}$ then $T_{40}(t) = (u_0)^{(8)}T_{41}(t)$ This is an important

consequence of the relation between $(\nu_1)^{(8)}$ and $(\bar{\nu}_1)^{(8)}$, **and definition of** $(u_0)^{(8)}$.

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**Proof** : From 99,20,44,22,23,44 we obtain

$$\frac{d\nu^{(9)}}{dt} = (a_{44})^{(9)} - \left((a'_{44})^{(9)} - (a'_{45})^{(9)} + (a''_{44})^{(9)}(T_{45},t)\right) - (a''_{45})^{(9)}(T_{45},t)\nu^{(9)} - (a_{45})^{(9)}\nu^{(9)}$$

**Definition of** $\nu^{(9)}$ :- $\boxed{\nu^{(9)} = \frac{G_{44}}{G_{45}}}$

It follows

$$-\left((a_{45})^{(9)}\left(\nu^{(9)}\right)^2 + (\sigma_2)^{(9)}\nu^{(9)} - (a_{44})^{(9)}\right) \le \frac{d\nu^{(9)}}{dt} \le -\left((a_{45})^{(9)}\left(\nu^{(9)}\right)^2 + (\sigma_1)^{(9)}\nu^{(9)} - (a_{44})^{(9)}\right)$$

From which one obtains

**Definition of** $(\bar{\nu}_1)^{(9)}, (\nu_0)^{(9)}$ :-

For $0 < \boxed{(\nu_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}}} < (\nu_1)^{(9)} < (\bar{\nu}_1)^{(9)}$

$$\nu^{(9)}(t) \ge \frac{(\nu_1)^{(9)} + (C)^{(9)}(\nu_2)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_0)^{(9)}\right)t\right]}}{1 + (C)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_0)^{(9)}\right)t\right]}} \quad , \quad \boxed{(C)^{(9)} = \frac{(\nu_1)^{(9)} - (\nu_0)^{(9)}}{(\nu_0)^{(9)} - (\nu_2)^{(9)}}}$$

it follows $(\nu_0)^{(9)} \le \nu^{(9)}(t) \le (\nu_9)^{(9)}$

In the same manner , we get

$$\nu^{(9)}(t) \le \frac{(\bar{\nu}_1)^{(9)} + (\bar{C})^{(9)}(\bar{\nu}_2)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}}{1 + (\bar{C})^{(9)} e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}} \quad , \quad \boxed{(\bar{C})^{(9)} = \frac{(\bar{\nu}_1)^{(9)} - (\nu_0)^{(9)}}{(\nu_0)^{(9)} - (\bar{\nu}_2)^{(9)}}}$$

From which we deduce $(\nu_0)^{(9)} \le \nu^{(9)}(t) \le (\bar{\nu}_1)^{(9)}$

If $0 < (\nu_1)^{(9)} < (\nu_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}} < (\bar{\nu}_1)^{(9)}$ we find like in the previous case,

$$(\nu_1)^{(9)} \le \frac{(\nu_1)^{(9)} + (C)^{(9)}(\nu_2)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_2)^{(9)}\right)t\right]}}{1 + (C)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\nu_1)^{(9)} - (\nu_2)^{(9)}\right)t\right]}} \le \nu^{(9)}(t) \le$$

$$\frac{(\bar{\nu}_1)^{(9)} + (\bar{C})^{(9)}(\bar{\nu}_2)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}}{1 + (\bar{C})^{(9)} e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}} \le (\bar{\nu}_1)^{(9)}$$

If $0 < (\nu_1)^{(9)} \le (\bar{\nu}_1)^{(9)} \le \boxed{(\nu_0)^{(9)} = \frac{G^0_{44}}{G^0_{45}}}$ , we obtain

$$(\nu_1)^{(9)} \le \nu^{(9)}(t) \le \frac{(\bar{\nu}_1)^{(9)} + (\bar{C})^{(9)}(\bar{\nu}_2)^{(9)} e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}}{1 + (\bar{C})^{(9)} e^{\left[-(a_{45})^{(9)}\left((\bar{\nu}_1)^{(9)} - (\bar{\nu}_2)^{(9)}\right)t\right]}} \le (\nu_0)^{(9)}$$

And so with the notation of the first part of condition (c) , we have

**Definition of** $\nu^{(9)}(t)$ :-

$(m_2)^{(9)} \leq \nu^{(9)}(t) \leq (m_1)^{(9)}$, $\boxed{\nu^{(9)}(t) = \frac{G_{44}(t)}{G_{45}(t)}}$

In a completely analogous way, we obtain

**Definition of** $u^{(9)}(t)$ :-

$(\mu_2)^{(9)} \leq u^{(9)}(t) \leq (\mu_1)^{(9)}$, $\boxed{u^{(9)}(t) = \frac{T_{44}(t)}{T_{45}(t)}}$

Now, using this result and replacing it in 99, 20,44,22,23, and 44 we get easily the result stated in the theorem.

**Particular case :**

If $(a_{44}^{''})^{(9)} = (a_{45}^{''})^{(9)}, then\ (\sigma_1)^{(9)} = (\sigma_2)^{(9)}$ and in this case $(\nu_1)^{(9)} = (\bar{\nu}_1)^{(9)}$ if in addition $(\nu_0)^{(9)} = (\nu_1)^{(9)}$ then $\nu^{(9)}(t) = (\nu_0)^{(9)}$ and as a consequence $G_{44}(t) = (\nu_0)^{(9)} G_{45}(t)$ **this also defines** $(\nu_0)^{(9)}$ **for the special case** .

Analogously if $(b_{44}^{''})^{(9)} = (b_{45}^{''})^{(9)}, then\ (\tau_1)^{(9)} = (\tau_2)^{(9)}$ and then $(u_1)^{(9)} = (\bar{u}_1)^{(9)}$if in addition $(u_0)^{(9)} = (u_1)^{(9)}$ then $T_{44}(t) = (u_0)^{(9)} T_{45}(t)$ This is an important consequence of the relation between $(\nu_1)^{(9)}$ and $(\bar{\nu}_1)^{(9)}$, **and definition of** $(u_0)^{(9)}$.

We can prove the following 425

**Theorem :** If $(a_i^{''})^{(1)} and\ (b_i^{''})^{(1)}$ are independent on $t$ , and the conditions with the notations

$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} < 0$

$(a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a_{13})^{(1)}(p_{13})^{(1)} + (a_{14}^{'})^{(1)}(p_{14})^{(1)} + (p_{13})^{(1)}(p_{14})^{(1)} > 0$

$(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} > 0$ ,

$(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)} - (b_{13})^{(1)}(b_{14})^{(1)} - (b_{13}^{'})^{(1)}(r_{14})^{(1)} - (b_{14}^{'})^{(1)}(r_{14})^{(1)} + (r_{13})^{(1)}(r_{14})^{(1)} < 0$

$with\ (p_{13})^{(1)}, (r_{14})^{(1)}$ as defined by equation are satisfied , then the system

**Theorem :** If $(a_i^{''})^{(2)} and\ (b_i^{''})^{(2)}$ are independent on t , and the conditions with the notations 426

$(a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} < 0$ 427

$(a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a_{16})^{(2)}(p_{16})^{(2)} + (a_{17}^{'})^{(2)}(p_{17})^{(2)} + (p_{16})^{(2)}(p_{17})^{(2)} > 0$ 428

$(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} > 0$ , 429

$(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)} - (b_{16})^{(2)}(b_{17})^{(2)} - (b_{16}^{'})^{(2)}(r_{17})^{(2)} - (b_{17}^{'})^{(2)}(r_{17})^{(2)} + (r_{16})^{(2)}(r_{17})^{(2)} < 0$ 430

$with\ (p_{16})^{(2)}, (r_{17})^{(2)}$ as defined by equation are satisfied , then the system

**Theorem :** If $(a_i^{''})^{(3)} and\ (b_i^{''})^{(3)}$ are independent on $t$ , and the conditions with the notations 431

$(a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} < 0$

| | |
|---|---|
| $(a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a_{20})^{(3)}(p_{20})^{(3)} + (a_{21}^{'})^{(3)}(p_{21})^{(3)} + (p_{20})^{(3)}(p_{21})^{(3)} > 0$<br>$(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} > 0$ ,<br>$(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)} - (b_{20})^{(3)}(b_{21})^{(3)} - (b_{20}^{'})^{(3)}(r_{21})^{(3)} - (b_{21}^{'})^{(3)}(r_{21})^{(3)} + (r_{20})^{(3)}(r_{21})^{(3)} < 0$<br>$with\ \ (p_{20})^{(3)}, (r_{21})^{(3)}$ as defined by equation are satisfied , then the system | |
| We can prove the following<br>**Theorem :** If $(a_i^{''})^{(4)} and\ (b_i^{''})^{(4)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} < 0$<br>$(a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a_{24})^{(4)}(p_{24})^{(4)} + (a_{25}^{'})^{(4)}(p_{25})^{(4)} + (p_{24})^{(4)}(p_{25})^{(4)} > 0$<br>$(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} > 0$ ,<br>$(b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} - (b_{24}^{'})^{(4)}(r_{25})^{(4)} - (b_{25}^{'})^{(4)}(r_{25})^{(4)} + (r_{24})^{(4)}(r_{25})^{(4)} < 0$<br>$with\ \ (p_{24})^{(4)}, (r_{25})^{(4)}$ as defined by equation are satisfied , then the system | 432 |
| **Theorem :** If $(a_i^{''})^{(5)} and\ (b_i^{''})^{(5)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} < 0$<br>$(a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a_{28})^{(5)}(p_{28})^{(5)} + (a_{29}^{'})^{(5)}(p_{29})^{(5)} + (p_{28})^{(5)}(p_{29})^{(5)} > 0$<br>$(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} > 0$ ,<br>$(b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} - (b_{28}^{'})^{(5)}(r_{29})^{(5)} - (b_{29}^{'})^{(5)}(r_{29})^{(5)} + (r_{28})^{(5)}(r_{29})^{(5)} < 0$<br>$with\ \ (p_{28})^{(5)}, (r_{29})^{(5)}$ as defined by equation are satisfied , then the system | 433 |
| **Theorem** If $(a_i^{''})^{(6)} and\ (b_i^{''})^{(6)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} < 0$<br>$(a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32})^{(6)}(p_{32})^{(6)} + (a_{33}^{'})^{(6)}(p_{33})^{(6)} + (p_{32})^{(6)}(p_{33})^{(6)} > 0$<br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} > 0$ ,<br>$(b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} - (b_{32}^{'})^{(6)}(r_{33})^{(6)} - (b_{33}^{'})^{(6)}(r_{33})^{(6)} + (r_{32})^{(6)}(r_{33})^{(6)} < 0$<br>$with\ \ (p_{32})^{(6)}, (r_{33})^{(6)}$ as defined by equation are satisfied , then the system | 434 |
| **Theorem :** If $(a_i^{''})^{(7)} and\ (b_i^{''})^{(7)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} < 0$ | 435 |

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| $(a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36})^{(7)}(p_{36})^{(7)} + (a_{37}^{'})^{(7)}(p_{37})^{(7)} + (p_{36})^{(7)}(p_{37})^{(7)} > 0$<br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} > 0$ ,<br>$(b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} - (b_{36}^{'})^{(7)}(r_{37})^{(7)} - (b_{37}^{'})^{(7)}(r_{37})^{(7)} + (r_{36})^{(7)}(r_{37})^{(7)} < 0$<br>$with\ \ (p_{36})^{(7)}, (r_{37})^{(7)}$ as defined by equation are satisfied , then the system | |
| **Theorem :** If $(a_i^{''})^{(8)} and\ (b_i^{''})^{(8)}$ are independent on $t$ , and the conditions with the notations<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} < 0$<br>$(a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40})^{(8)}(p_{40})^{(8)} + (a_{41}^{'})^{(8)}(p_{41})^{(8)} + (p_{40})^{(8)}(p_{41})^{(8)} > 0$<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} > 0$ ,<br>$(b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} - (b_{40}^{'})^{(8)}(r_{41})^{(8)} - (b_{41}^{'})^{(8)}(r_{41})^{(8)} + (r_{40})^{(8)}(r_{41})^{(8)} < 0$<br>$with\ \ (p_{40})^{(8)}, (r_{41})^{(8)}$ as defined by equation are satisfied , then the system | 436 |
| **Theorem :** If $(a_i^{''})^{(9)} and\ (b_i^{''})^{(9)}$ are independent on $t$ , and the conditions (with the notations 45,46,27,28)<br>$(a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} < 0$<br>$(a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44})^{(9)}(p_{44})^{(9)} + (a_{45}^{'})^{(9)}(p_{45})^{(9)} + (p_{44})^{(9)}(p_{45})^{(9)} > 0$<br>$(b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} > 0$ ,<br>$(b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} - (b_{44}^{'})^{(9)}(r_{45})^{(9)} - (b_{45}^{'})^{(9)}(r_{45})^{(9)} + (r_{44})^{(9)}(r_{45})^{(9)} < 0$<br>$with\ \ (p_{44})^{(9)}, (r_{45})^{(9)}$ as defined by equation 45 are satisfied , then the system | 436 A |
| $(a_{13})^{(1)}G_{14} - \left[(a_{13}^{'})^{(1)} + (a_{13}^{''})^{(1)}(T_{14})\right]G_{13} = 0$ | 437 |
| $(a_{14})^{(1)}G_{13} - \left[(a_{14}^{'})^{(1)} + (a_{14}^{''})^{(1)}(T_{14})\right]G_{14} = 0$ | 438 |
| $(a_{15})^{(1)}G_{14} - \left[(a_{15}^{'})^{(1)} + (a_{15}^{''})^{(1)}(T_{14})\right]G_{15} = 0$ | 439 |
| $(b_{13})^{(1)}T_{14} - [(b_{13}^{'})^{(1)} - (b_{13}^{''})^{(1)}(G)\ ]T_{13} = 0$ | 440 |
| $(b_{14})^{(1)}T_{13} - [(b_{14}^{'})^{(1)} - (b_{14}^{''})^{(1)}(G)\ ]T_{14} = 0$ | 441 |
| $(b_{15})^{(1)}T_{14} - [(b_{15}^{'})^{(1)} - (b_{15}^{''})^{(1)}(G)\ ]T_{15} = 0$ | 442 |
| has a unique positive solution , which is an equilibrium solution for the system | |

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| $(a_{16})^{(2)}G_{17} - \left[(a_{16}^{'})^{(2)} + (a_{16}^{''})^{(2)}(T_{17})\right]G_{16} = 0$ | 443 |
| $(a_{17})^{(2)}G_{16} - \left[(a_{17}^{'})^{(2)} + (a_{17}^{''})^{(2)}(T_{17})\right]G_{17} = 0$ | 444 |
| $(a_{18})^{(2)}G_{17} - \left[(a_{18}^{'})^{(2)} + (a_{18}^{''})^{(2)}(T_{17})\right]G_{18} = 0$ | 445 |
| $(b_{16})^{(2)}T_{17} - [(b_{16}^{'})^{(2)} - (b_{16}^{''})^{(2)}(G_{19})\,]T_{16} = 0$ | 446 |
| $(b_{17})^{(2)}T_{16} - [(b_{17}^{'})^{(2)} - (b_{17}^{''})^{(2)}(G_{19})\,]T_{17} = 0$ | 447 |
| $(b_{18})^{(2)}T_{17} - [(b_{18}^{'})^{(2)} - (b_{18}^{''})^{(2)}(G_{19})\,]T_{18} = 0$ | 448 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{20})^{(3)}G_{21} - \left[(a_{20}^{'})^{(3)} + (a_{20}^{''})^{(3)}(T_{21})\right]G_{20} = 0$ | 449 |
| $(a_{21})^{(3)}G_{20} - \left[(a_{21}^{'})^{(3)} + (a_{21}^{''})^{(3)}(T_{21})\right]G_{21} = 0$ | 450 |
| $(a_{22})^{(3)}G_{21} - \left[(a_{22}^{'})^{(3)} + (a_{22}^{''})^{(3)}(T_{21})\right]G_{22} = 0$ | 451 |
| $(b_{20})^{(3)}T_{21} - [(b_{20}^{'})^{(3)} - (b_{20}^{''})^{(3)}(G_{23})\,]T_{20} = 0$ | 452 |
| $(b_{21})^{(3)}T_{20} - [(b_{21}^{'})^{(3)} - (b_{21}^{''})^{(3)}(G_{23})\,]T_{21} = 0$ | 453 |
| $(b_{22})^{(3)}T_{21} - [(b_{22}^{'})^{(3)} - (b_{22}^{''})^{(3)}(G_{23})\,]T_{22} = 0$ | 454 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{24})^{(4)}G_{25} - \left[(a_{24}^{'})^{(4)} + (a_{24}^{''})^{(4)}(T_{25})\right]G_{24} = 0$ | 455 |
| $(a_{25})^{(4)}G_{24} - \left[(a_{25}^{'})^{(4)} + (a_{25}^{''})^{(4)}(T_{25})\right]G_{25} = 0$ | 456 |
| $(a_{26})^{(4)}G_{25} - \left[(a_{26}^{'})^{(4)} + (a_{26}^{''})^{(4)}(T_{25})\right]G_{26} = 0$ | 457 |
| $(b_{24})^{(4)}T_{25} - [(b_{24}^{'})^{(4)} - (b_{24}^{''})^{(4)}\left((G_{27})\right)\,]T_{24} = 0$ | 458 |
| $(b_{25})^{(4)}T_{24} - [(b_{25}^{'})^{(4)} - (b_{25}^{''})^{(4)}\left((G_{27})\right)\,]T_{25} = 0$ | 459 |
| $(b_{26})^{(4)}T_{25} - [(b_{26}^{'})^{(4)} - (b_{26}^{''})^{(4)}\left((G_{27})\right)\,]T_{26} = 0$ | 460 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{28})^{(5)}G_{29} - \left[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29})\right]G_{28} = 0$ | 461 |
| $(a_{29})^{(5)}G_{28} - \left[(a_{29}^{'})^{(5)} + (a_{29}^{''})^{(5)}(T_{29})\right]G_{29} = 0$ | 462 |

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| $(a_{30})^{(5)}G_{29} - \left[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29})\right]G_{30} = 0$ | 463 |
| $(b_{28})^{(5)}T_{29} - [(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}(G_{31})\,]T_{28} = 0$ | 464 |
| $(b_{29})^{(5)}T_{28} - [(b_{29}^{'})^{(5)} - (b_{29}^{''})^{(5)}(G_{31})\,]T_{29} = 0$ | 465 |
| $(b_{30})^{(5)}T_{29} - [(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}(G_{31})\,]T_{30} = 0$ | 466 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{32})^{(6)}G_{33} - \left[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33})\right]G_{32} = 0$ | 467 |
| $(a_{33})^{(6)}G_{32} - \left[(a_{33}^{'})^{(6)} + (a_{33}^{''})^{(6)}(T_{33})\right]G_{33} = 0$ | 468 |
| $(a_{34})^{(6)}G_{33} - \left[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33})\right]G_{34} = 0$ | 469 |
| $(b_{32})^{(6)}T_{33} - [(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}(G_{35})\,]T_{32} = 0$ | 470 |
| $(b_{33})^{(6)}T_{32} - [(b_{33}^{'})^{(6)} - (b_{33}^{''})^{(6)}(G_{35})\,]T_{33} = 0$ | 471 |
| $(b_{34})^{(6)}T_{33} - [(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}(G_{35})\,]T_{34} = 0$ | 472 |
| has a unique positive solution , which is an equilibrium solution | |
| $(a_{36})^{(7)}G_{37} - \left[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37})\right]G_{36} = 0$ | 473 |
| $(a_{37})^{(7)}G_{36} - \left[(a_{37}^{'})^{(7)} + (a_{37}^{''})^{(7)}(T_{37})\right]G_{37} = 0$ | 474 |
| $(a_{38})^{(7)}G_{37} - \left[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37})\right]G_{38} = 0$ | 475 |
| $(b_{36})^{(7)}T_{37} - [(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}(G_{39})]T_{36} = 0$ | 476 |
| $(b_{37})^{(7)}T_{36} - [(b_{37}^{'})^{(7)} - (b_{37}^{''})^{(7)}(G_{39})\,]T_{37} = 0$ | 477 |
| $(b_{38})^{(7)}T_{37} - [(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}(G_{39})\,]T_{38} = 0$ | 478 |
| $(a_{40})^{(8)}G_{41} - \left[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41})\right]G_{40} = 0$ | 479 |
| $(a_{41})^{(8)}G_{40} - \left[(a_{41}^{'})^{(8)} + (a_{41}^{''})^{(8)}(T_{41})\right]G_{41} = 0$ | 480 |

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| $(a_{42})^{(8)}G_{41} - \left[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41})\right]G_{42} = 0$ | 481 |
| $(b_{40})^{(8)}T_{41} - [(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}(G_{43})\,]T_{40} = 0$ | 482 |
| $(b_{41})^{(8)}T_{40} - [(b_{41}^{'})^{(8)} - (b_{41}^{''})^{(8)}(G_{43})\,]T_{41} = 0$ | 483 |
| $(b_{42})^{(8)}T_{41} - [(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}(G_{43})\,]T_{42} = 0$ | 484 |
| $(a_{44})^{(9)}G_{45} - \left[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45})\right]G_{44} = 0$ | 484 A |
| $(a_{45})^{(9)}G_{44} - \left[(a_{45}^{'})^{(9)} + (a_{45}^{''})^{(9)}(T_{45})\right]G_{45} = 0$ | |
| $(a_{46})^{(9)}G_{45} - \left[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45})\right]G_{46} = 0$ | |
| $(b_{44})^{(9)}T_{45} - [(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}(G_{47})\,]T_{44} = 0$ | |
| $(b_{45})^{(9)}T_{44} - [(b_{45}^{'})^{(9)} - (b_{45}^{''})^{(9)}(G_{47})\,]T_{45} = 0$ | |
| $(b_{46})^{(9)}T_{45} - [(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}(G_{47})\,]T_{46} = 0$ | |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{13}, G_{14}$ if<br><br>$F(T) = (a_{13}^{'})^{(1)}(a_{14}^{'})^{(1)} - (a_{13})^{(1)}(a_{14})^{(1)} + (a_{13}^{'})^{(1)}(a_{14}^{''})^{(1)}(T_{14}) + (a_{14}^{'})^{(1)}(a_{13}^{''})^{(1)}(T_{14}) + (a_{13}^{''})^{(1)}(T_{14})(a_{14}^{''})^{(1)}(T_{14}) = 0$ | 485 |
| **Proof:**<br><br>(h) Indeed the first two equations have a nontrivial solution $G_{16}, G_{17}$ if<br><br>$\mathrm{F}(T_{19}) = (a_{16}^{'})^{(2)}(a_{17}^{'})^{(2)} - (a_{16})^{(2)}(a_{17})^{(2)} + (a_{16}^{'})^{(2)}(a_{17}^{''})^{(2)}(T_{17}) + (a_{17}^{'})^{(2)}(a_{16}^{''})^{(2)}(T_{17}) + (a_{16}^{''})^{(2)}(T_{17})(a_{17}^{''})^{(2)}(T_{17}) = 0$ | 486 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{20}, G_{21}$ if<br><br>$F(T_{23}) = (a_{20}^{'})^{(3)}(a_{21}^{'})^{(3)} - (a_{20})^{(3)}(a_{21})^{(3)} + (a_{20}^{'})^{(3)}(a_{21}^{''})^{(3)}(T_{21}) + (a_{21}^{'})^{(3)}(a_{20}^{''})^{(3)}(T_{21}) + (a_{20}^{''})^{(3)}(T_{21})(a_{21}^{''})^{(3)}(T_{21}) = 0$ | 487 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{24}, G_{25}$ if<br><br>$F(T_{27}) = (a_{24}^{'})^{(4)}(a_{25}^{'})^{(4)} - (a_{24})^{(4)}(a_{25})^{(4)} + (a_{24}^{'})^{(4)}(a_{25}^{''})^{(4)}(T_{25}) + (a_{25}^{'})^{(4)}(a_{24}^{''})^{(4)}(T_{25}) +$ | 488 |

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| $(a_{24}^{''})^{(4)}(T_{25})(a_{25}^{''})^{(4)}(T_{25}) = 0$ | |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{28}, G_{29}$ if<br><br>$F(T_{31}) = (a_{28}^{'})^{(5)}(a_{29}^{'})^{(5)} - (a_{28})^{(5)}(a_{29})^{(5)} + (a_{28}^{'})^{(5)}(a_{29}^{''})^{(5)}(T_{29}) + (a_{29}^{'})^{(5)}(a_{28}^{''})^{(5)}(T_{29}) + (a_{28}^{''})^{(5)}(T_{29})(a_{29}^{''})^{(5)}(T_{29}) = 0$ | 489 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{32}, G_{33}$ if<br><br>$F(T_{35}) = (a_{32}^{'})^{(6)}(a_{33}^{'})^{(6)} - (a_{32})^{(6)}(a_{33})^{(6)} + (a_{32}^{'})^{(6)}(a_{33}^{''})^{(6)}(T_{33}) + (a_{33}^{'})^{(6)}(a_{32}^{''})^{(6)}(T_{33}) + (a_{32}^{''})^{(6)}(T_{33})(a_{33}^{''})^{(6)}(T_{33}) = 0$ | 490 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{36}, G_{37}$ if<br><br>$F(T_{39}) = (a_{36}^{'})^{(7)}(a_{37}^{'})^{(7)} - (a_{36})^{(7)}(a_{37})^{(7)} + (a_{36}^{'})^{(7)}(a_{37}^{''})^{(7)}(T_{37}) + (a_{37}^{'})^{(7)}(a_{36}^{''})^{(7)}(T_{37}) + (a_{36}^{''})^{(7)}(T_{37})(a_{37}^{''})^{(7)}(T_{37}) = 0$ | 491 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{40}, G_{41}$ if<br><br>$F(T_{43}) = (a_{40}^{'})^{(8)}(a_{41}^{'})^{(8)} - (a_{40})^{(8)}(a_{41})^{(8)} + (a_{40}^{'})^{(8)}(a_{41}^{''})^{(8)}(T_{41}) + (a_{41}^{'})^{(8)}(a_{40}^{''})^{(8)}(T_{41}) + (a_{40}^{''})^{(8)}(T_{41})(a_{41}^{''})^{(8)}(T_{41}) = 0$ | 492 |
| **Proof:**<br><br>(a) Indeed the first two equations have a nontrivial solution $G_{44}, G_{45}$ if<br><br>$F(T_{47}) = (a_{44}^{'})^{(9)}(a_{45}^{'})^{(9)} - (a_{44})^{(9)}(a_{45})^{(9)} + (a_{44}^{'})^{(9)}(a_{45}^{''})^{(9)}(T_{45}) + (a_{45}^{'})^{(9)}(a_{44}^{''})^{(9)}(T_{45}) + (a_{44}^{''})^{(9)}(T_{45})(a_{45}^{''})^{(9)}(T_{45}) = 0$ | 492 A |
| **Definition and uniqueness of** $T_{14}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{14})$ being increasing, it follows that there exists a unique $T_{14}^{*}$ for which $f(T_{14}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{13} = \frac{(a_{13})^{(1)}G_{14}}{[(a_{13}^{'})^{(1)}+(a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15} = \frac{(a_{15})^{(1)}G_{14}}{[(a_{15}^{'})^{(1)}+(a_{15}^{''})^{(1)}(T_{14}^{*})]}$ | 493 |
| **Definition and uniqueness of** $T_{17}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(2)}(T_{17})$ being increasing, it follows that there exists a unique $T_{17}^{*}$ for which $f(T_{17}^{*}) = 0$. With this value , we obtain from the three first | 494 |

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| equations | |
| $G_{16} = \frac{(a_{16})^{(2)}G_{17}}{[(a_{16}^{'})^{(2)}+(a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18} = \frac{(a_{18})^{(2)}G_{17}}{[(a_{18}^{'})^{(2)}+(a_{18}^{''})^{(2)}(T_{17}^{*})]}$ | 495 |
| **Definition and uniqueness of** $T_{21}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(1)}(T_{21})$ being increasing, it follows that there exists a unique $T_{21}^{*}$ for which $f(T_{21}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{20} = \frac{(a_{20})^{(3)}G_{21}}{[(a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21}^{*})]}$ , $G_{22} = \frac{(a_{22})^{(3)}G_{21}}{[(a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21}^{*})]}$ | 496 |
| **Definition and uniqueness of** $T_{25}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(4)}(T_{25})$ being increasing, it follows that there exists a unique $T_{25}^{*}$ for which $f(T_{25}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{24} = \frac{(a_{24})^{(4)}G_{25}}{[(a_{24}^{'})^{(4)}+(a_{24}^{''})^{(4)}(T_{25}^{*})]}$ , $G_{26} = \frac{(a_{26})^{(4)}G_{25}}{[(a_{26}^{'})^{(4)}+(a_{26}^{''})^{(4)}(T_{25}^{*})]}$ | 497 |
| **Definition and uniqueness of** $T_{29}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(5)}(T_{29})$ being increasing, it follows that there exists a unique $T_{29}^{*}$ for which $f(T_{29}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{28} = \frac{(a_{28})^{(5)}G_{29}}{[(a_{28}^{'})^{(5)}+(a_{28}^{''})^{(5)}(T_{29}^{*})]}$ , $G_{30} = \frac{(a_{30})^{(5)}G_{29}}{[(a_{30}^{'})^{(5)}+(a_{30}^{''})^{(5)}(T_{29}^{*})]}$ | 498 |
| **Definition and uniqueness of** $T_{33}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(6)}(T_{33})$ being increasing, it follows that there exists a unique $T_{33}^{*}$ for which $f(T_{33}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{32} = \frac{(a_{32})^{(6)}G_{33}}{[(a_{32}^{'})^{(6)}+(a_{32}^{''})^{(6)}(T_{33}^{*})]}$ , $G_{34} = \frac{(a_{34})^{(6)}G_{33}}{[(a_{34}^{'})^{(6)}+(a_{34}^{''})^{(6)}(T_{33}^{*})]}$ | 499 |
| **Definition and uniqueness of** $T_{37}^{*}$ :-<br><br>After hypothesis $f(0) < 0, f(\infty) > 0$ and the functions $(a_i^{''})^{(7)}(T_{37})$ being increasing, it follows that there exists a unique $T_{37}^{*}$ for which $f(T_{37}^{*}) = 0$. With this value , we obtain from the three first equations<br><br>$G_{36} = \frac{(a_{36})^{(7)}G_{37}}{[(a_{36}^{'})^{(7)}+(a_{36}^{''})^{(7)}(T_{37}^{*})]}$ , $G_{38} = \frac{(a_{38})^{(7)}G_{37}}{[(a_{38}^{'})^{(7)}+(a_{38}^{''})^{(7)}(T_{37}^{*})]}$ | 500 |

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| **Definition and uniqueness of** $T_{41}^{*}$ :-<br><br>After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(8)}(T_{41})$ being increasing, it follows that there exists a unique $T_{41}^{*}$ for which $f(T_{41}^{*})=0$. With this value , we obtain from the three first equations<br><br>$G_{40}=\frac{(a_{40})^{(8)}G_{41}}{[(a_{40}^{'})^{(8)}+(a_{40}^{''})^{(8)}(T_{41}^{*})]}$ , $G_{42}=\frac{(a_{42})^{(8)}G_{41}}{[(a_{42}^{'})^{(8)}+(a_{42}^{''})^{(8)}(T_{41}^{*})]}$ | 501 |
| **Definition and uniqueness of** $T_{45}^{*}$ :-<br><br>After hypothesis $f(0)<0, f(\infty)>0$ and the functions $(a_i^{''})^{(9)}(T_{45})$ being increasing, it follows that there exists a unique $T_{45}^{*}$ for which $f(T_{45}^{*})=0$. With this value , we obtain from the three first equations<br><br>$G_{44}=\frac{(a_{44})^{(9)}G_{45}}{[(a_{44}^{'})^{(9)}+(a_{44}^{''})^{(9)}(T_{45}^{*})]}$ , $G_{46}=\frac{(a_{46})^{(9)}G_{45}}{[(a_{46}^{'})^{(9)}+(a_{46}^{''})^{(9)}(T_{45}^{*})]}$ | 501 A |
| By the same argument, the equations admit solutions $G_{13}, G_{14}$ if<br>$\varphi(G)=(b_{13}^{'})^{(1)}(b_{14}^{'})^{(1)}-(b_{13})^{(1)}(b_{14})^{(1)}-$<br>$[(b_{13}^{'})^{(1)}(b_{14}^{''})^{(1)}(G)+(b_{14}^{'})^{(1)}(b_{13}^{''})^{(1)}(G)]+(b_{13}^{''})^{(1)}(G)(b_{14}^{''})^{(1)}(G)=0$<br><br>Where in $G(G_{13},G_{14},G_{15}), G_{13}, G_{15}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{14}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{14}^{*}$ such that $\varphi(G^{*})=0$ | 502 |
| By the same argument, the equations admit solutions $G_{16}, G_{17}$ if<br><br>$\varphi(G_{19})=(b_{16}^{'})^{(2)}(b_{17}^{'})^{(2)}-(b_{16})^{(2)}(b_{17})^{(2)}-$<br>$[(b_{16}^{'})^{(2)}(b_{17}^{''})^{(2)}(G_{19})+(b_{17}^{'})^{(2)}(b_{16}^{''})^{(2)}(G_{19})]+(b_{16}^{''})^{(2)}(G_{19})(b_{17}^{''})^{(2)}(G_{19})=0$ | 503 |
| Where in $(G_{19})(G_{16},G_{17},G_{18}), G_{16}, G_{18}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{17}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{14}^{*}$ such that $\varphi((G_{19})^{*})=0$ | 504 |
| By the same argument, the equations admit solutions $G_{20}, G_{21}$ if<br><br>$\varphi(G_{23})=(b_{20}^{'})^{(3)}(b_{21}^{'})^{(3)}-(b_{20})^{(3)}(b_{21})^{(3)}-$<br>$[(b_{20}^{'})^{(3)}(b_{21}^{''})^{(3)}(G_{23})+(b_{21}^{'})^{(3)}(b_{20}^{''})^{(3)}(G_{23})]+(b_{20}^{''})^{(3)}(G_{23})(b_{21}^{''})^{(3)}(G_{23})=0$<br><br>Where in $G_{23}(G_{20},G_{21},G_{22}), G_{20}, G_{22}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{21}$ taking into account the hypothesis $\varphi(0)>0$ , $\varphi(\infty)<0$ it follows that there exists a unique $G_{21}^{*}$ such that $\varphi((G_{23})^{*})=0$ | 505 |

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| By the same argument, the equations admit solutions $G_{24}, G_{25}$ if<br><br>$\varphi(G_{27}) = (b_{24}^{'})^{(4)}(b_{25}^{'})^{(4)} - (b_{24})^{(4)}(b_{25})^{(4)} -$<br><br>$\left[(b_{24}^{'})^{(4)}(b_{25}^{''})^{(4)}(G_{27}) + (b_{25}^{'})^{(4)}(b_{24}^{''})^{(4)}(G_{27})\right] + (b_{24}^{''})^{(4)}(G_{27})(b_{25}^{''})^{(4)}(G_{27}) = 0$<br><br>Where in $(G_{27})(G_{24}, G_{25}, G_{26}), G_{24}, G_{26}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{25}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{25}^{*}$ such that $\varphi((G_{27})^{*}) = 0$ | 506 |
| By the same argument, the equations admit solutions $G_{28}, G_{29}$ if<br>$\varphi(G_{31}) = (b_{28}^{'})^{(5)}(b_{29}^{'})^{(5)} - (b_{28})^{(5)}(b_{29})^{(5)} -$<br><br>$\left[(b_{28}^{'})^{(5)}(b_{29}^{''})^{(5)}(G_{31}) + (b_{29}^{'})^{(5)}(b_{28}^{''})^{(5)}(G_{31})\right] + (b_{28}^{''})^{(5)}(G_{31})(b_{29}^{''})^{(5)}(G_{31}) = 0$<br><br>Where in $(G_{31})(G_{28}, G_{29}, G_{30}), G_{28}, G_{30}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{29}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{29}^{*}$ such that $\varphi((G_{31})^{*}) = 0$ | 507 |
| By the same argument, the equations admit solutions $G_{32}, G_{33}$ if<br>$\varphi(G_{35}) = (b_{32}^{'})^{(6)}(b_{33}^{'})^{(6)} - (b_{32})^{(6)}(b_{33})^{(6)} -$<br><br>$\left[(b_{32}^{'})^{(6)}(b_{33}^{''})^{(6)}(G_{35}) + (b_{33}^{'})^{(6)}(b_{32}^{''})^{(6)}(G_{35})\right] + (b_{32}^{''})^{(6)}(G_{35})(b_{33}^{''})^{(6)}(G_{35}) = 0$<br><br>Where in $(G_{35})(G_{32}, G_{33}, G_{34}), G_{32}, G_{34}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{33}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{33}^{*}$ such that $\varphi({G_{35}}^{*}) = 0$ | 508 |
| By the same argument, the equations admit solutions $G_{36}, G_{37}$ if<br>$\varphi(G_{39}) = (b_{36}^{'})^{(7)}(b_{37}^{'})^{(7)} - (b_{36})^{(7)}(b_{37})^{(7)} -$<br>$\left[(b_{36}^{'})^{(7)}(b_{37}^{''})^{(7)}(G_{39}) + (b_{37}^{'})^{(7)}(b_{36}^{''})^{(7)}(G_{39})\right] + (b_{36}^{''})^{(7)}(G_{39})(b_{37}^{''})^{(7)}(G_{39}) = 0$<br><br>Where in $(G_{39})(G_{36}, G_{37}, G_{38}), G_{36}, G_{38}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{37}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{37}^{*}$ such that $\varphi({G_{39}}^{*}) = 0$ | 509 |
| By the same argument, the equations admit solutions $G_{40}, G_{41}$ if<br>$\varphi(G_{43}) = (b_{40}^{'})^{(8)}(b_{41}^{'})^{(8)} - (b_{40})^{(8)}(b_{41})^{(8)} -$<br>$\left[(b_{40}^{'})^{(8)}(b_{41}^{''})^{(8)}(G_{43}) + (b_{41}^{'})^{(8)}(b_{40}^{''})^{(8)}(G_{43})\right] + (b_{40}^{''})^{(8)}(G_{43})(b_{41}^{''})^{(8)}(G_{43}) = 0$<br><br>Where in $(G_{43})(G_{40}, G_{41}, G_{42}), G_{40}, G_{42}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{41}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{41}^{*}$ such that $\varphi({G_{43}}^{*}) = 0$ | 510 |
| By the same argument, the equations 92,93 admit solutions $G_{44}, G_{45}$ if<br>$\varphi(G_{47}) = (b_{44}^{'})^{(9)}(b_{45}^{'})^{(9)} - (b_{44})^{(9)}(b_{45})^{(9)} -$ | |

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| $\left[(b_{44}^{'})^{(9)}(b_{45}^{''})^{(9)}(G_{47}) + (b_{45}^{'})^{(9)}(b_{44}^{''})^{(9)}(G_{47})\right]+(b_{44}^{''})^{(9)}(G_{47})(b_{45}^{''})^{(9)}(G_{47}) = 0$<br><br>Where in $(G_{47})(G_{44}, G_{45}, G_{46}), G_{44}, G_{46}$ must be replaced by their values from 96. It is easy to see that φ is a decreasing function in $G_{45}$ taking into account the hypothesis $\varphi(0) > 0$ , $\varphi(\infty) < 0$ it follows that there exists a unique $G_{45}^{*}$ such that $\varphi((G_{47})^{*}) = 0$ | |
| Finally we obtain the unique solution<br><br>$G_{14}^{*}$ given by $\varphi(G^{*}) = 0$ , $T_{14}^{*}$ given by $f(T_{14}^{*}) = 0$ and<br><br>$G_{13}^{*} = \frac{(a_{13})^{(1)}G_{14}^{*}}{[(a_{13}^{'})^{(1)}+(a_{13}^{''})^{(1)}(T_{14}^{*})]}$ , $G_{15}^{*} = \frac{(a_{15})^{(1)}G_{14}^{*}}{[(a_{15}^{'})^{(1)}+(a_{15}^{''})^{(1)}(T_{14}^{*})]}$<br><br>$T_{13}^{*} = \frac{(b_{13})^{(1)}T_{14}^{*}}{[(b_{13}^{'})^{(1)}-(b_{13}^{''})^{(1)}(G^{*})]}$ , $T_{15}^{*} = \frac{(b_{15})^{(1)}T_{14}^{*}}{[(b_{15}^{'})^{(1)}-(b_{15}^{''})^{(1)}(G^{*})]}$<br><br>Obviously, these values represent an equilibrium solution | 511 |
| Finally we obtain the unique solution | |
| $G_{17}^{*}$ given by $\varphi((G_{19})^{*}) = 0$ , $T_{17}^{*}$ given by $f(T_{17}^{*}) = 0$ and | 512 |
| $G_{16}^{*} = \frac{(a_{16})^{(2)}G_{17}^{*}}{[(a_{16}^{'})^{(2)}+(a_{16}^{''})^{(2)}(T_{17}^{*})]}$ , $G_{18}^{*} = \frac{(a_{18})^{(2)}G_{17}^{*}}{[(a_{18}^{'})^{(2)}+(a_{18}^{''})^{(2)}(T_{17}^{*})]}$ | 513 |
| $T_{16}^{*} = \frac{(b_{16})^{(2)}T_{17}^{*}}{[(b_{16}^{'})^{(2)}-(b_{16}^{''})^{(2)}((G_{19})^{*})]}$ , $T_{18}^{*} = \frac{(b_{18})^{(2)}T_{17}^{*}}{[(b_{18}^{'})^{(2)}-(b_{18}^{''})^{(2)}((G_{19})^{*})]}$ | 514 |
| Obviously, these values represent an equilibrium solution | |
| Finally we obtain the unique solution<br><br>$G_{21}^{*}$ given by $\varphi((G_{23})^{*}) = 0$ , $T_{21}^{*}$ given by $f(T_{21}^{*}) = 0$ and<br><br>$G_{20}^{*} = \frac{(a_{20})^{(3)}G_{21}^{*}}{[(a_{20}^{'})^{(3)}+(a_{20}^{''})^{(3)}(T_{21}^{*})]}$ , $G_{22}^{*} = \frac{(a_{22})^{(3)}G_{21}^{*}}{[(a_{22}^{'})^{(3)}+(a_{22}^{''})^{(3)}(T_{21}^{*})]}$<br><br>$T_{20}^{*} = \frac{(b_{20})^{(3)}T_{21}^{*}}{[(b_{20}^{'})^{(3)}-(b_{20}^{''})^{(3)}(G_{23}{}^{*})]}$ , $T_{22}^{*} = \frac{(b_{22})^{(3)}T_{21}^{*}}{[(b_{22}^{'})^{(3)}-(b_{22}^{''})^{(3)}(G_{23}{}^{*})]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 515 |
| Finally we obtain the unique solution<br><br>$G_{25}^{*}$ given by $\varphi(G_{27}) = 0$ , $T_{25}^{*}$ given by $f(T_{25}^{*}) = 0$ and<br><br>$G_{24}^{*} = \frac{(a_{24})^{(4)}G_{25}^{*}}{[(a_{24}^{'})^{(4)}+(a_{24}^{''})^{(4)}(T_{25}^{*})]}$ , $G_{26}^{*} = \frac{(a_{26})^{(4)}G_{25}^{*}}{[(a_{26}^{'})^{(4)}+(a_{26}^{''})^{(4)}(T_{25}^{*})]}$ | 516 |
| $T_{24}^{*} = \frac{(b_{24})^{(4)}T_{25}^{*}}{[(b_{24}^{'})^{(4)}-(b_{24}^{''})^{(4)}((G_{27})^{*})]}$ , $T_{26}^{*} = \frac{(b_{26})^{(4)}T_{25}^{*}}{[(b_{26}^{'})^{(4)}-(b_{26}^{''})^{(4)}((G_{27})^{*})]}$ | 517 |

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| Obviously, these values represent an equilibrium solution of global equations | |
| Finally we obtain the unique solution<br><br>$G_{29}^{*}$ given by $\varphi((G_{31})^{*}) = 0$ , $T_{29}^{*}$ given by $f(T_{29}^{*}) = 0$ and<br><br>$G_{28}^{*} = \frac{(a_{28})^{(5)} G_{29}^{*}}{[(a_{28}^{'})^{(5)} + (a_{28}^{''})^{(5)}(T_{29}^{*})]}$ , $G_{30}^{*} = \frac{(a_{30})^{(5)} G_{29}^{*}}{[(a_{30}^{'})^{(5)} + (a_{30}^{''})^{(5)}(T_{29}^{*})]}$ | 518 |
| $T_{28}^{*} = \frac{(b_{28})^{(5)} T_{29}^{*}}{[(b_{28}^{'})^{(5)} - (b_{28}^{''})^{(5)}((G_{31})^{*})]}$ , $T_{30}^{*} = \frac{(b_{30})^{(5)} T_{29}^{*}}{[(b_{30}^{'})^{(5)} - (b_{30}^{''})^{(5)}((G_{31})^{*})]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 519 |
| Finally we obtain the unique solution<br><br>$G_{33}^{*}$ given by $\varphi((G_{35})^{*}) = 0$ , $T_{33}^{*}$ given by $f(T_{33}^{*}) = 0$ and<br><br>$G_{32}^{*} = \frac{(a_{32})^{(6)} G_{33}^{*}}{[(a_{32}^{'})^{(6)} + (a_{32}^{''})^{(6)}(T_{33}^{*})]}$ , $G_{34}^{*} = \frac{(a_{34})^{(6)} G_{33}^{*}}{[(a_{34}^{'})^{(6)} + (a_{34}^{''})^{(6)}(T_{33}^{*})]}$ | 520 |
| $T_{32}^{*} = \frac{(b_{32})^{(6)} T_{33}^{*}}{[(b_{32}^{'})^{(6)} - (b_{32}^{''})^{(6)}((G_{35})^{*})]}$ , $T_{34}^{*} = \frac{(b_{34})^{(6)} T_{33}^{*}}{[(b_{34}^{'})^{(6)} - (b_{34}^{''})^{(6)}((G_{35})^{*})]}$<br><br>Obviously, these values represent an equilibrium solution of global equations | 521 |
| Finally we obtain the unique solution<br><br>$G_{37}^{*}$ given by $\varphi((G_{39})^{*}) = 0$ , $T_{37}^{*}$ given by $f(T_{37}^{*}) = 0$ and<br><br>$G_{36}^{*} = \frac{(a_{36})^{(7)} G_{37}^{*}}{[(a_{36}^{'})^{(7)} + (a_{36}^{''})^{(7)}(T_{37}^{*})]}$ , $G_{38}^{*} = \frac{(a_{38})^{(7)} G_{37}^{*}}{[(a_{38}^{'})^{(7)} + (a_{38}^{''})^{(7)}(T_{37}^{*})]}$<br><br>$T_{36}^{*} = \frac{(b_{36})^{(7)} T_{37}^{*}}{[(b_{36}^{'})^{(7)} - (b_{36}^{''})^{(7)}((G_{39})^{*})]}$ , $T_{38}^{*} = \frac{(b_{38})^{(7)} T_{37}^{*}}{[(b_{38}^{'})^{(7)} - (b_{38}^{''})^{(7)}((G_{39})^{*})]}$ | 522 |
| Finally we obtain the unique solution<br><br>$G_{41}^{*}$ given by $\varphi((G_{43})^{*}) = 0$ , $T_{41}^{*}$ given by $f(T_{41}^{*}) = 0$ and<br><br>$G_{40}^{*} = \frac{(a_{40})^{(8)} G_{41}^{*}}{[(a_{40}^{'})^{(8)} + (a_{40}^{''})^{(8)}(T_{41}^{*})]}$ , $G_{42}^{*} = \frac{(a_{42})^{(8)} G_{41}^{*}}{[(a_{42}^{'})^{(8)} + (a_{42}^{''})^{(8)}(T_{41}^{*})]}$<br><br>$T_{40}^{*} = \frac{(b_{40})^{(8)} T_{41}^{*}}{[(b_{40}^{'})^{(8)} - (b_{40}^{''})^{(8)}((G_{43})^{*})]}$ , $T_{42}^{*} = \frac{(b_{42})^{(8)} T_{41}^{*}}{[(b_{42}^{'})^{(8)} - (b_{42}^{''})^{(8)}((G_{43})^{*})]}$ | 523 |
| Finally we obtain the unique solution of 89 to 99 | 523 A |

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| $G_{45}^{*}$ given by $\varphi((G_{47})^{*}) = 0$ , $T_{45}^{*}$ given by $f(T_{45}^{*}) = 0$ and<br>$G_{44}^{*} = \frac{(a_{44})^{(9)} G_{45}^{*}}{[(a_{44}^{'})^{(9)} + (a_{44}^{''})^{(9)}(T_{45}^{*})]}$ , $G_{46}^{*} = \frac{(a_{46})^{(9)} G_{45}^{*}}{[(a_{46}^{'})^{(9)} + (a_{46}^{''})^{(9)}(T_{45}^{*})]}$<br>$T_{44}^{*} = \frac{(b_{44})^{(9)} T_{45}^{*}}{[(b_{44}^{'})^{(9)} - (b_{44}^{''})^{(9)}((G_{47})^{*})]}$ , $T_{46}^{*} = \frac{(b_{46})^{(9)} T_{45}^{*}}{[(b_{46}^{'})^{(9)} - (b_{46}^{''})^{(9)}((G_{47})^{*})]}$ | |
| **ASYMPTOTIC STABILITY ANALYSIS**<br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(1)}$ *and* $(b_i^{''})^{(1)}$ Belong to $C^{(1)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br>**Proof:** Denote<br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br>$G_i = G_i^{*} + \mathbb{G}_i$ , $T_i = T_i^{*} + \mathbb{T}_i$<br>$\frac{\partial (a_{14}^{''})^{(1)}}{\partial T_{14}}(T_{14}^{*}) = (q_{14})^{(1)}$ , $\frac{\partial (b_i^{''})^{(1)}}{\partial G_j}(G^{*}) = s_{ij}$ | 524 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{13}}{dt} = -\left((a_{13}^{'})^{(1)} + (p_{13})^{(1)}\right)\mathbb{G}_{13} + (a_{13})^{(1)}\mathbb{G}_{14} - (q_{13})^{(1)} G_{13}^{*}\mathbb{T}_{14}$ | 525 |
| $\frac{d\mathbb{G}_{14}}{dt} = -\left((a_{14}^{'})^{(1)} + (p_{14})^{(1)}\right)\mathbb{G}_{14} + (a_{14})^{(1)}\mathbb{G}_{13} - (q_{14})^{(1)} G_{14}^{*}\mathbb{T}_{14}$ | 526 |
| $\frac{d\mathbb{G}_{15}}{dt} = -\left((a_{15}^{'})^{(1)} + (p_{15})^{(1)}\right)\mathbb{G}_{15} + (a_{15})^{(1)}\mathbb{G}_{14} - (q_{15})^{(1)} G_{15}^{*}\mathbb{T}_{14}$ | 527 |
| $\frac{d\mathbb{T}_{13}}{dt} = -\left((b_{13}^{'})^{(1)} - (r_{13})^{(1)}\right)\mathbb{T}_{13} + (b_{13})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(13)(j)} T_{13}^{*}\mathbb{G}_j\right)$ | 528 |
| $\frac{d\mathbb{T}_{14}}{dt} = -\left((b_{14}^{'})^{(1)} - (r_{14})^{(1)}\right)\mathbb{T}_{14} + (b_{14})^{(1)}\mathbb{T}_{13} + \sum_{j=13}^{15}\left(s_{(14)(j)} T_{14}^{*}\mathbb{G}_j\right)$ | 529 |
| $\frac{d\mathbb{T}_{15}}{dt} = -\left((b_{15}^{'})^{(1)} - (r_{15})^{(1)}\right)\mathbb{T}_{15} + (b_{15})^{(1)}\mathbb{T}_{14} + \sum_{j=13}^{15}\left(s_{(15)(j)} T_{15}^{*}\mathbb{G}_j\right)$ | 530 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(2)}$ and $(b_i^{''})^{(2)}$ Belong to $C^{(2)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable | 531 |
| Proof: Denote<br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- | |
| $G_i = G_i^{*} + \mathbb{G}_i$ , $T_i = T_i^{*} + \mathbb{T}_i$ | 532 |

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| $\frac{\partial(a_{17}^{''})^{(2)}}{\partial \mathrm{T}_{17}}(\mathrm{T}_{17}^{*}) = (q_{17})^{(2)}$ , $\frac{\partial(b_i^{''})^{(2)}}{\partial \mathrm{G}_j}((G_{19})^{*}) = s_{ij}$ | 533 |
| taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{\mathrm{d}\mathbb{G}_{16}}{\mathrm{dt}} = -\left((a_{16}^{'})^{(2)} + (p_{16})^{(2)}\right)\mathbb{G}_{16} + (a_{16})^{(2)}\mathbb{G}_{17} - (q_{16})^{(2)}\mathrm{G}_{16}^{*}\mathbb{T}_{17}$ | 534 |
| $\frac{\mathrm{d}\mathbb{G}_{17}}{\mathrm{dt}} = -\left((a_{17}^{'})^{(2)} + (p_{17})^{(2)}\right)\mathbb{G}_{17} + (a_{17})^{(2)}\mathbb{G}_{16} - (q_{17})^{(2)}\mathrm{G}_{17}^{*}\mathbb{T}_{17}$ | 535 |
| $\frac{\mathrm{d}\mathbb{G}_{18}}{\mathrm{dt}} = -\left((a_{18}^{'})^{(2)} + (p_{18})^{(2)}\right)\mathbb{G}_{18} + (a_{18})^{(2)}\mathbb{G}_{17} - (q_{18})^{(2)}\mathrm{G}_{18}^{*}\mathbb{T}_{17}$ | 536 |
| $\frac{\mathrm{d}\mathbb{T}_{16}}{\mathrm{dt}} = -\left((b_{16}^{'})^{(2)} - (r_{16})^{(2)}\right)\mathbb{T}_{16} + (b_{16})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(16)(j)}\mathrm{T}_{16}^{*}\mathbb{G}_j\right)$ | 537 |
| $\frac{\mathrm{d}\mathbb{T}_{17}}{\mathrm{dt}} = -\left((b_{17}^{'})^{(2)} - (r_{17})^{(2)}\right)\mathbb{T}_{17} + (b_{17})^{(2)}\mathbb{T}_{16} + \sum_{j=16}^{18}\left(s_{(17)(j)}\mathrm{T}_{17}^{*}\mathbb{G}_j\right)$ | 538 |
| $\frac{\mathrm{d}\mathbb{T}_{18}}{\mathrm{dt}} = -\left((b_{18}^{'})^{(2)} - (r_{18})^{(2)}\right)\mathbb{T}_{18} + (b_{18})^{(2)}\mathbb{T}_{17} + \sum_{j=16}^{18}\left(s_{(18)(j)}\mathrm{T}_{18}^{*}\mathbb{G}_j\right)$ | 539 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(3)}$ *and* $(b_i^{''})^{(3)}$ Belong to $C^{(3)}(\mathbb{R}_{+})$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote<br><br>**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^{*} + \mathbb{G}_i \quad , T_i = T_i^{*} + \mathbb{T}_i$<br><br>$\frac{\partial(a_{21}^{''})^{(3)}}{\partial T_{21}}(T_{21}^{*}) = (q_{21})^{(3)}$ , $\frac{\partial(b_i^{''})^{(3)}}{\partial G_j}((G_{23})^{*}) = s_{ij}$ | 540 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{20}}{dt} = -\left((a_{20}^{'})^{(3)} + (p_{20})^{(3)}\right)\mathbb{G}_{20} + (a_{20})^{(3)}\mathbb{G}_{21} - (q_{20})^{(3)}G_{20}^{*}\mathbb{T}_{21}$ | 541 |
| $\frac{d\mathbb{G}_{21}}{dt} = -\left((a_{21}^{'})^{(3)} + (p_{21})^{(3)}\right)\mathbb{G}_{21} + (a_{21})^{(3)}\mathbb{G}_{20} - (q_{21})^{(3)}G_{21}^{*}\mathbb{T}_{21}$ | 542 |
| $\frac{d\mathbb{G}_{22}}{dt} = -\left((a_{22}^{'})^{(3)} + (p_{22})^{(3)}\right)\mathbb{G}_{22} + (a_{22})^{(3)}\mathbb{G}_{21} - (q_{22})^{(3)}G_{22}^{*}\mathbb{T}_{21}$ | 543 |
| $\frac{d\mathbb{T}_{20}}{dt} = -\left((b_{20}^{'})^{(3)} - (r_{20})^{(3)}\right)\mathbb{T}_{20} + (b_{20})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(20)(j)}T_{20}^{*}\mathbb{G}_j\right)$ | 544 |
| $\frac{d\mathbb{T}_{21}}{dt} = -\left((b_{21}^{'})^{(3)} - (r_{21})^{(3)}\right)\mathbb{T}_{21} + (b_{21})^{(3)}\mathbb{T}_{20} + \sum_{j=20}^{22}\left(s_{(21)(j)}T_{21}^{*}\mathbb{G}_j\right)$ | 545 |
| $\frac{d\mathbb{T}_{22}}{dt} = -\left((b_{22}^{'})^{(3)} - (r_{22})^{(3)}\right)\mathbb{T}_{22} + (b_{22})^{(3)}\mathbb{T}_{21} + \sum_{j=20}^{22}\left(s_{(22)(j)}T_{22}^{*}\mathbb{G}_j\right)$ | 546 |
| **ASYMPTOTIC STABILITY ANALYSIS** | 547 |

**Theorem 4:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(4)}$ $and$ $(b_i^{''})^{(4)}$ Belong to $C^{(4)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.

Proof: Denote

**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-

$$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$$

$$\frac{\partial (a_{25}^{''})^{(4)}}{\partial T_{25}}(T_{25}^*) = (q_{25})^{(4)} \quad , \frac{\partial (b_i^{''})^{(4)}}{\partial G_j}((G_{27})^* ) = s_{ij}$$

Then taking into account equations and neglecting the terms of power 2, we obtain

$$\frac{d\mathbb{G}_{24}}{dt} = -\left((a_{24}^{'})^{(4)} + (p_{24})^{(4)}\right)\mathbb{G}_{24} + (a_{24})^{(4)}\mathbb{G}_{25} - (q_{24})^{(4)}G_{24}^*\mathbb{T}_{25} \quad (549)$$

$$\frac{d\mathbb{G}_{25}}{dt} = -\left((a_{25}^{'})^{(4)} + (p_{25})^{(4)}\right)\mathbb{G}_{25} + (a_{25})^{(4)}\mathbb{G}_{24} - (q_{25})^{(4)}G_{25}^*\mathbb{T}_{25} \quad (550)$$

$$\frac{d\mathbb{G}_{26}}{dt} = -\left((a_{26}^{'})^{(4)} + (p_{26})^{(4)}\right)\mathbb{G}_{26} + (a_{26})^{(4)}\mathbb{G}_{25} - (q_{26})^{(4)}G_{26}^*\mathbb{T}_{25} \quad (551)$$

$$\frac{d\mathbb{T}_{24}}{dt} = -\left((b_{24}^{'})^{(4)} - (r_{24})^{(4)}\right)\mathbb{T}_{24} + (b_{24})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(24)(j)}T_{24}^*\mathbb{G}_j\right) \quad (552)$$

$$\frac{d\mathbb{T}_{25}}{dt} = -\left((b_{25}^{'})^{(4)} - (r_{25})^{(4)}\right)\mathbb{T}_{25} + (b_{25})^{(4)}\mathbb{T}_{24} + \sum_{j=24}^{26}\left(s_{(25)(j)}T_{25}^*\mathbb{G}_j\right) \quad (553)$$

$$\frac{d\mathbb{T}_{26}}{dt} = -\left((b_{26}^{'})^{(4)} - (r_{26})^{(4)}\right)\mathbb{T}_{26} + (b_{26})^{(4)}\mathbb{T}_{25} + \sum_{j=24}^{26}\left(s_{(26)(j)}T_{26}^*\mathbb{G}_j\right) \quad (554)$$

**ASYMPTOTIC STABILITY ANALYSIS** (555)

**Theorem 5:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(5)}$ $and$ $(b_i^{''})^{(5)}$ Belong to $C^{(5)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.

Proof: Denote

**Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :- (556)

$$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$$

$$\frac{\partial (a_{29}^{''})^{(5)}}{\partial T_{29}}(T_{29}^*) = (q_{29})^{(5)} \quad , \frac{\partial (b_i^{''})^{(5)}}{\partial G_j}(\,(G_{31})^*\,) = s_{ij}$$

Then taking into account equations and neglecting the terms of power 2, we obtain

$$\frac{d\mathbb{G}_{28}}{dt} = -\left((a_{28}^{'})^{(5)} + (p_{28})^{(5)}\right)\mathbb{G}_{28} + (a_{28})^{(5)}\mathbb{G}_{29} - (q_{28})^{(5)}G_{28}^*\mathbb{T}_{29} \quad (557)$$

$$\frac{d\mathbb{G}_{29}}{dt} = -\left((a_{29}^{'})^{(5)} + (p_{29})^{(5)}\right)\mathbb{G}_{29} + (a_{29})^{(5)}\mathbb{G}_{28} - (q_{29})^{(5)}G_{29}^*\mathbb{T}_{29} \quad (558)$$

$$\frac{d\mathbb{G}_{30}}{dt} = -\left((a_{30}^{'})^{(5)} + (p_{30})^{(5)}\right)\mathbb{G}_{30} + (a_{30})^{(5)}\mathbb{G}_{29} - (q_{30})^{(5)}G_{30}^*\mathbb{T}_{29} \quad (559)$$

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| $\frac{d\mathbb{T}_{28}}{dt} = -\left((b_{28}^{'})^{(5)} - (r_{28})^{(5)}\right)\mathbb{T}_{28} + (b_{28})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(28)(j)}T_{28}^{*}\mathbb{G}_{j}\right)$ | 560 |
| $\frac{d\mathbb{T}_{29}}{dt} = -\left((b_{29}^{'})^{(5)} - (r_{29})^{(5)}\right)\mathbb{T}_{29} + (b_{29})^{(5)}\mathbb{T}_{28} + \sum_{j=28}^{30}\left(s_{(29)(j)}T_{29}^{*}\mathbb{G}_{j}\right)$ | 561 |
| $\frac{d\mathbb{T}_{30}}{dt} = -\left((b_{30}^{'})^{(5)} - (r_{30})^{(5)}\right)\mathbb{T}_{30} + (b_{30})^{(5)}\mathbb{T}_{29} + \sum_{j=28}^{30}\left(s_{(30)(j)}T_{30}^{*}\mathbb{G}_{j}\right)$ | 562 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 6:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(6)}$ $and$ $(b_i^{''})^{(6)}$ Belong to $C^{(6)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 563 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial(a_{33}^{''})^{(6)}}{\partial T_{33}}(T_{33}^{*}) = (q_{33})^{(6)}$ , $\frac{\partial(b_i^{''})^{(6)}}{\partial G_j}((G_{35})^{*}) = s_{ij}$ | 564 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{32}}{dt} = -\left((a_{32}^{'})^{(6)} + (p_{32})^{(6)}\right)\mathbb{G}_{32} + (a_{32})^{(6)}\mathbb{G}_{33} - (q_{32})^{(6)}G_{32}^{*}\mathbb{T}_{33}$ | 565 |
| $\frac{d\mathbb{G}_{33}}{dt} = -\left((a_{33}^{'})^{(6)} + (p_{33})^{(6)}\right)\mathbb{G}_{33} + (a_{33})^{(6)}\mathbb{G}_{32} - (q_{33})^{(6)}G_{33}^{*}\mathbb{T}_{33}$ | 566 |
| $\frac{d\mathbb{G}_{34}}{dt} = -\left((a_{34}^{'})^{(6)} + (p_{34})^{(6)}\right)\mathbb{G}_{34} + (a_{34})^{(6)}\mathbb{G}_{33} - (q_{34})^{(6)}G_{34}^{*}\mathbb{T}_{33}$ | 567 |
| $\frac{d\mathbb{T}_{32}}{dt} = -\left((b_{32}^{'})^{(6)} - (r_{32})^{(6)}\right)\mathbb{T}_{32} + (b_{32})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(32)(j)}T_{32}^{*}\mathbb{G}_{j}\right)$ | 568 |
| $\frac{d\mathbb{T}_{33}}{dt} = -\left((b_{33}^{'})^{(6)} - (r_{33})^{(6)}\right)\mathbb{T}_{33} + (b_{33})^{(6)}\mathbb{T}_{32} + \sum_{j=32}^{34}\left(s_{(33)(j)}T_{33}^{*}\mathbb{G}_{j}\right)$ | 569 |
| $\frac{d\mathbb{T}_{34}}{dt} = -\left((b_{34}^{'})^{(6)} - (r_{34})^{(6)}\right)\mathbb{T}_{34} + (b_{34})^{(6)}\mathbb{T}_{33} + \sum_{j=32}^{34}\left(s_{(34)(j)}T_{34}^{*}\mathbb{G}_{j}\right)$ | 570 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 7:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(7)}$ $and$ $(b_i^{''})^{(7)}$ Belong to $C^{(7)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | 571 |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$ | 572 |

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|---|---|
| $\frac{\partial (a_{37}^{''})^{(7)}}{\partial T_{37}}(T_{37}^{*}) = (q_{37})^{(7)}$ , $\frac{\partial (b_i^{''})^{(7)}}{\partial G_j}\left((G_{39})^{**}\right) = s_{ij}$ | |
| Then taking into account equations and neglecting the terms of power 2, we obtain from | |
| $\frac{d\mathbb{G}_{36}}{dt} = -\left((a_{36}^{'})^{(7)} + (p_{36})^{(7)}\right)\mathbb{G}_{36} + (a_{36})^{(7)}\mathbb{G}_{37} - (q_{36})^{(7)}G_{36}^{*}\mathbb{T}_{37}$ | 573 |
| $\frac{d\mathbb{G}_{37}}{dt} = -\left((a_{37}^{'})^{(7)} + (p_{37})^{(7)}\right)\mathbb{G}_{37} + (a_{37})^{(7)}\mathbb{G}_{36} - (q_{37})^{(7)}G_{37}^{*}\mathbb{T}_{37}$ | 574 |
| $\frac{d\mathbb{G}_{38}}{dt} = -\left((a_{38}^{'})^{(7)} + (p_{38})^{(7)}\right)\mathbb{G}_{38} + (a_{38})^{(7)}\mathbb{G}_{37} - (q_{38})^{(7)}G_{38}^{*}\mathbb{T}_{37}$ | 575 |
| $\frac{d\mathbb{T}_{36}}{dt} = -\left((b_{36}^{'})^{(7)} - (r_{36})^{(7)}\right)\mathbb{T}_{36} + (b_{36})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(36)(j)}T_{36}^{*}\mathbb{G}_j\right)$ | 576 |
| $\frac{d\mathbb{T}_{37}}{dt} = -\left((b_{37}^{'})^{(7)} - (r_{37})^{(7)}\right)\mathbb{T}_{37} + (b_{37})^{(7)}\mathbb{T}_{36} + \sum_{j=36}^{38}\left(s_{(37)(j)}T_{37}^{*}\mathbb{G}_j\right)$ | 578 |
| $\frac{d\mathbb{T}_{38}}{dt} = -\left((b_{38}^{'})^{(7)} - (r_{38})^{(7)}\right)\mathbb{T}_{38} + (b_{38})^{(7)}\mathbb{T}_{37} + \sum_{j=36}^{38}\left(s_{(38)(j)}T_{38}^{*}\mathbb{G}_j\right)$ | 579 |
| Obviously, these values represent an equilibrium solution<br><br>**ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 8:** If the conditions of the previous theorem are satisfied and if the functions $(a_i^{''})^{(8)}$ *and* $(b_i^{''})^{(8)}$ Belong to $C^{(8)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br><br>Proof: Denote | |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^{*} + \mathbb{G}_i \quad , T_i = T_i^{*} + \mathbb{T}_i$<br><br>$\frac{\partial (a_{41}^{''})^{(8)}}{\partial T_{41}}(T_{41}^{*}) = (q_{41})^{(8)}$ , $\frac{\partial (b_i^{''})^{(8)}}{\partial G_j}\left((G_{43})^{*}\right) = s_{ij}$ | 580 |
| Then taking into account equations and neglecting the terms of power 2, we obtain | |
| $\frac{d\mathbb{G}_{40}}{dt} = -\left((a_{40}^{'})^{(8)} + (p_{40})^{(8)}\right)\mathbb{G}_{40} + (a_{40})^{(8)}\mathbb{G}_{41} - (q_{40})^{(8)}G_{40}^{*}\mathbb{T}_{41}$ | 581 |
| $\frac{d\mathbb{G}_{41}}{dt} = -\left((a_{41}^{'})^{(8)} + (p_{41})^{(8)}\right)\mathbb{G}_{41} + (a_{41})^{(8)}\mathbb{G}_{40} - (q_{41})^{(8)}G_{41}^{*}\mathbb{T}_{41}$ | 582 |
| $\frac{d\mathbb{G}_{42}}{dt} = -\left((a_{42}^{'})^{(8)} + (p_{42})^{(8)}\right)\mathbb{G}_{42} + (a_{42})^{(8)}\mathbb{G}_{41} - (q_{42})^{(8)}G_{42}^{*}\mathbb{T}_{41}$ | 583 |
| $\frac{d\mathbb{T}_{40}}{dt} = -\left((b_{40}^{'})^{(8)} - (r_{40})^{(8)}\right)\mathbb{T}_{40} + (b_{40})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(40)(j)}T_{40}^{*}\mathbb{G}_j\right)$ | 584 |

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|---|---|
| $\frac{d\mathbb{T}_{41}}{dt} = -\left((b'_{41})^{(8)} - (r_{41})^{(8)}\right)\mathbb{T}_{41} + (b_{41})^{(8)}\mathbb{T}_{40} + \sum_{j=40}^{42}\left(s_{(41)(j)}T^*_{41}\mathbb{G}_j\right)$ | 585 |
| $\frac{d\mathbb{T}_{42}}{dt} = -\left((b'_{42})^{(8)} - (r_{42})^{(8)}\right)\mathbb{T}_{42} + (b_{42})^{(8)}\mathbb{T}_{41} + \sum_{j=40}^{42}\left(s_{(42)(j)}T^*_{42}\mathbb{G}_j\right)$ | 586 |
| **ASYMPTOTIC STABILITY ANALYSIS**<br><br>**Theorem 9:** If the conditions of the previous theorem are satisfied and if the functions $(a''_i)^{(9)}$ *and* $(b''_i)^{(9)}$ Belong to $C^{(9)}(\mathbb{R}_+)$ then the above equilibrium point is asymptotically stable.<br>**Proof: Denote** | 586 A |
| **Definition of** $\mathbb{G}_i, \mathbb{T}_i$ :-<br><br>$G_i = G_i^* + \mathbb{G}_i \quad , T_i = T_i^* + \mathbb{T}_i$<br><br>$\frac{\partial(a''_{45})^{(9)}}{\partial T_{45}}(T^*_{45}) = (q_{45})^{(9)} \ , \ \frac{\partial(b''_i)^{(9)}}{\partial G_j}((G_{47})^*) = s_{ij}$ | |
| Then taking into account equations 89 to 99 and neglecting the terms of power 2, we obtain from 99 to 44 | |
| $\frac{d\mathbb{G}_{44}}{dt} = -\left((a'_{44})^{(9)} + (p_{44})^{(9)}\right)\mathbb{G}_{44} + (a_{44})^{(9)}\mathbb{G}_{45} - (q_{44})^{(9)}G^*_{44}\mathbb{T}_{45}$ | 586 B |
| $\frac{d\mathbb{G}_{45}}{dt} = -\left((a'_{45})^{(9)} + (p_{45})^{(9)}\right)\mathbb{G}_{45} + (a_{45})^{(9)}\mathbb{G}_{44} - (q_{45})^{(9)}G^*_{45}\mathbb{T}_{45}$ | 586 C |
| $\frac{d\mathbb{G}_{46}}{dt} = -\left((a'_{46})^{(9)} + (p_{46})^{(9)}\right)\mathbb{G}_{46} + (a_{46})^{(9)}\mathbb{G}_{45} - (q_{46})^{(9)}G^*_{46}\mathbb{T}_{45}$ | 586 D |
| $\frac{d\mathbb{T}_{44}}{dt} = -\left((b'_{44})^{(9)} - (r_{44})^{(9)}\right)\mathbb{T}_{44} + (b_{44})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(44)(j)}T^*_{44}\mathbb{G}_j\right)$ | 586 E |
| $\frac{d\mathbb{T}_{45}}{dt} = -\left((b'_{45})^{(9)} - (r_{45})^{(9)}\right)\mathbb{T}_{45} + (b_{45})^{(9)}\mathbb{T}_{44} + \sum_{j=44}^{46}\left(s_{(45)(j)}T^*_{45}\mathbb{G}_j\right)$ | 586 F |
| $\frac{d\mathbb{T}_{46}}{dt} = -\left((b'_{46})^{(9)} - (r_{46})^{(9)}\right)\mathbb{T}_{46} + (b_{46})^{(9)}\mathbb{T}_{45} + \sum_{j=44}^{46}\left(s_{(46)(j)}T^*_{46}\mathbb{G}_j\right)$ | 586 G |
| **The characteristic equation of this system is** | 587 |
| $\left((\lambda)^{(1)} + (b'_{15})^{(1)} - (r_{15})^{(1)}\right)\{\left((\lambda)^{(1)} + (a'_{15})^{(1)} + (p_{15})^{(1)}\right)$<br><br>$\left[\left((\lambda)^{(1)} + (a'_{13})^{(1)} + (p_{13})^{(1)}\right)(q_{14})^{(1)}G^*_{14} + (a_{14})^{(1)}(q_{13})^{(1)}G^*_{13}\right)\Big]$<br><br>$\left((\lambda)^{(1)} + (b'_{13})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(14)}T^*_{14} + (b_{14})^{(1)}s_{(13),(14)}T^*_{14}\Big)$<br><br>$+\left((\lambda)^{(1)} + (a'_{14})^{(1)} + (p_{14})^{(1)}\right)(q_{13})^{(1)}G^*_{13} + (a_{13})^{(1)}(q_{14})^{(1)}G^*_{14}\Big)$<br><br>$\left((\lambda)^{(1)} + (b'_{13})^{(1)} - (r_{13})^{(1)}\right)s_{(14),(13)}T^*_{14} + (b_{14})^{(1)}s_{(13),(13)}T^*_{13}\Big)$<br><br>$\left((\lambda)^{(1)}\right)^2 + \left((a'_{13})^{(1)} + (a'_{14})^{(1)} + (p_{13})^{(1)} + (p_{14})^{(1)}\right)(\lambda)^{(1)}\Big)$ | |

$$\left((\lambda)^{(1)}\right)^2+\left((b'_{13})^{(1)}+(b'_{14})^{(1)}-(r_{13})^{(1)}+(r_{14})^{(1)}\right)(\lambda)^{(1)}\Big)$$

$$+\left(\left((\lambda)^{(1)}\right)^2+\left((a'_{13})^{(1)}+(a'_{14})^{(1)}+(p_{13})^{(1)}+(p_{14})^{(1)}\right)(\lambda)^{(1)}\right)(q_{15})^{(1)}G_{15}$$

$$+\left((\lambda)^{(1)}+(a'_{13})^{(1)}+(p_{13})^{(1)}\right)\left((a_{15})^{(1)}(q_{14})^{(1)}G^*_{14}+(a_{14})^{(1)}(a_{15})^{(1)}(q_{13})^{(1)}G^*_{13}\right)$$

$$\left((\lambda)^{(1)}+(b'_{13})^{(1)}-(r_{13})^{(1)}\right)s_{(14),(15)}T^*_{14}+(b_{14})^{(1)}s_{(13),(15)}T^*_{13}\Big)\}=0$$

$$+$$

$$(\lambda)^{(2)}+(b'_{18})^{(2)}-(r_{18})^{(2)}\}\{\left((\lambda)^{(2)}+(a'_{18})^{(2)}+(p_{18})^{(2)}\right)$$

$$\left[\left((\lambda)^{(2)}+(a'_{16})^{(2)}+(p_{16})^{(2)}\right)(q_{17})^{(2)}\mathrm{G}^*_{17}+(a_{17})^{(2)}(q_{16})^{(2)}\mathrm{G}^*_{16}\right)\Big]$$

$$\left((\lambda)^{(2)}+(b'_{16})^{(2)}-(r_{16})^{(2)}\right)s_{(17),(17)}\mathrm{T}^*_{17}+(b_{17})^{(2)}s_{(16),(17)}\mathrm{T}^*_{17}\Big)$$

$$+\left(\left((\lambda)^{(2)}+(a'_{17})^{(2)}+(p_{17})^{(2)}\right)(q_{16})^{(2)}\mathrm{G}^*_{16}+(a_{16})^{(2)}(q_{17})^{(2)}\mathrm{G}^*_{17}\right)$$

$$\left((\lambda)^{(2)}+(b'_{16})^{(2)}-(r_{16})^{(2)}\right)s_{(17),(16)}\mathrm{T}^*_{17}+(b_{17})^{(2)}s_{(16),(16)}\mathrm{T}^*_{16}\Big)$$

$$\left((\lambda)^{(2)}\right)^2+\left((a'_{16})^{(2)}+(a'_{17})^{(2)}+(p_{16})^{(2)}+(p_{17})^{(2)}\right)(\lambda)^{(2)}\Big)$$

$$\left((\lambda)^{(2)}\right)^2+\left((b'_{16})^{(2)}+(b'_{17})^{(2)}-(r_{16})^{(2)}+(r_{17})^{(2)}\right)(\lambda)^{(2)}\Big)$$

$$+\left(\left((\lambda)^{(2)}\right)^2+\left((a'_{16})^{(2)}+(a'_{17})^{(2)}+(p_{16})^{(2)}+(p_{17})^{(2)}\right)(\lambda)^{(2)}\right)(q_{18})^{(2)}\mathrm{G}_{18}$$

$$+\left((\lambda)^{(2)}+(a'_{16})^{(2)}+(p_{16})^{(2)}\right)\left((a_{18})^{(2)}(q_{17})^{(2)}\mathrm{G}^*_{17}+(a_{17})^{(2)}(a_{18})^{(2)}(q_{16})^{(2)}\mathrm{G}^*_{16}\right)$$

$$\left((\lambda)^{(2)}+(b'_{16})^{(2)}-(r_{16})^{(2)}\right)s_{(17),(18)}\mathrm{T}^*_{17}+(b_{17})^{(2)}s_{(16),(18)}\mathrm{T}^*_{16}\Big)\}=0$$

$$+$$

$$(\lambda)^{(3)}+(b'_{22})^{(3)}-(r_{22})^{(3)}\}\{\left((\lambda)^{(3)}+(a'_{22})^{(3)}+(p_{22})^{(3)}\right)$$

$$\left[\left((\lambda)^{(3)}+(a'_{20})^{(3)}+(p_{20})^{(3)}\right)(q_{21})^{(3)}G^*_{21}+(a_{21})^{(3)}(q_{20})^{(3)}G^*_{20}\right)\Big]$$

$$\left((\lambda)^{(3)}+(b'_{20})^{(3)}-(r_{20})^{(3)}\right)s_{(21),(21)}T^*_{21}+(b_{21})^{(3)}s_{(20),(21)}T^*_{21}\Big)$$

$$+\left(\left((\lambda)^{(3)}+(a'_{21})^{(3)}+(p_{21})^{(3)}\right)(q_{20})^{(3)}G^*_{20}+(a_{20})^{(3)}(q_{21})^{(1)}G^*_{21}\right)$$

$$\left((\lambda)^{(3)}+(b'_{20})^{(3)}-(r_{20})^{(3)}\right)s_{(21),(20)}T^*_{21}+(b_{21})^{(3)}s_{(20),(20)}T^*_{20}\Big)$$

$$\left((\lambda)^{(3)}\right)^2+\left((a'_{20})^{(3)}+(a'_{21})^{(3)}+(p_{20})^{(3)}+(p_{21})^{(3)}\right)(\lambda)^{(3)}\Big)$$

$$\left(\left((\lambda)^{(3)}\right)^2+\left((b_{20}')^{(3)}+(b_{21}')^{(3)}-(r_{20})^{(3)}+(r_{21})^{(3)}\right)(\lambda)^{(3)}\right)$$

$$+\left(\left((\lambda)^{(3)}\right)^2+\left((a_{20}')^{(3)}+(a_{21}')^{(3)}+(p_{20})^{(3)}+(p_{21})^{(3)}\right)(\lambda)^{(3)}\right)(q_{22})^{(3)}G_{22}$$

$$+\left((\lambda)^{(3)}+(a_{20}')^{(3)}+(p_{20})^{(3)}\right)\left((a_{22})^{(3)}(q_{21})^{(3)}G_{21}^{*}+(a_{21})^{(3)}(a_{22})^{(3)}(q_{20})^{(3)}G_{20}^{*}\right)$$

$$\left(\left((\lambda)^{(3)}+(b_{20}')^{(3)}-(r_{20})^{(3)}\right)s_{(21),(22)}T_{21}^{*}+(b_{21})^{(3)}s_{(20),(22)}T_{20}^{*}\right)\}=0$$

$$+$$

$$(\lambda)^{(4)}+(b_{26}')^{(4)}-(r_{26})^{(4)}\{\left((\lambda)^{(4)}+(a_{26}')^{(4)}+(p_{26})^{(4)}\right)$$

$$\left[\left(\left((\lambda)^{(4)}+(a_{24}')^{(4)}+(p_{24})^{(4)}\right)(q_{25})^{(4)}G_{25}^{*}+(a_{25})^{(4)}(q_{24})^{(4)}G_{24}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(4)}+(b_{24}')^{(4)}-(r_{24})^{(4)}\right)s_{(25),(25)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(25)}T_{25}^{*}\right)$$

$$+\left(\left((\lambda)^{(4)}+(a_{25}')^{(4)}+(p_{25})^{(4)}\right)(q_{24})^{(4)}G_{24}^{*}+(a_{24})^{(4)}(q_{25})^{(4)}G_{25}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}+(b_{24}')^{(4)}-(r_{24})^{(4)}\right)s_{(25),(24)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(24)}T_{24}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}\right)^2+\left((a_{24}')^{(4)}+(a_{25}')^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)$$

$$\left(\left((\lambda)^{(4)}\right)^2+\left((b_{24}')^{(4)}+(b_{25}')^{(4)}-(r_{24})^{(4)}+(r_{25})^{(4)}\right)(\lambda)^{(4)}\right)$$

$$+\left(\left((\lambda)^{(4)}\right)^2+\left((a_{24}')^{(4)}+(a_{25}')^{(4)}+(p_{24})^{(4)}+(p_{25})^{(4)}\right)(\lambda)^{(4)}\right)(q_{26})^{(4)}G_{26}$$

$$+\left((\lambda)^{(4)}+(a_{24}')^{(4)}+(p_{24})^{(4)}\right)\left((a_{26})^{(4)}(q_{25})^{(4)}G_{25}^{*}+(a_{25})^{(4)}(a_{26})^{(4)}(q_{24})^{(4)}G_{24}^{*}\right)$$

$$\left(\left((\lambda)^{(4)}+(b_{24}')^{(4)}-(r_{24})^{(4)}\right)s_{(25),(26)}T_{25}^{*}+(b_{25})^{(4)}s_{(24),(26)}T_{24}^{*}\right)\}=0$$

$$+$$

$$(\lambda)^{(5)}+(b_{30}')^{(5)}-(r_{30})^{(5)}\{\left((\lambda)^{(5)}+(a_{30}')^{(5)}+(p_{30})^{(5)}\right)$$

$$\left[\left(\left((\lambda)^{(5)}+(a_{28}')^{(5)}+(p_{28})^{(5)}\right)(q_{29})^{(5)}G_{29}^{*}+(a_{29})^{(5)}(q_{28})^{(5)}G_{28}^{*}\right)\right]$$

$$\left(\left((\lambda)^{(5)}+(b_{28}')^{(5)}-(r_{28})^{(5)}\right)s_{(29),(29)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(29)}T_{29}^{*}\right)$$

$$+\left(\left((\lambda)^{(5)}+(a_{29}')^{(5)}+(p_{29})^{(5)}\right)(q_{28})^{(5)}G_{28}^{*}+(a_{28})^{(5)}(q_{29})^{(5)}G_{29}^{*}\right)$$

$$\left(\left((\lambda)^{(5)}+(b_{28}')^{(5)}-(r_{28})^{(5)}\right)s_{(29),(28)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(28)}T_{28}^{*}\right)$$

$$\left(\left((\lambda)^{(5)}\right)^2+\left((a_{28}')^{(5)}+(a_{29}')^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)$$

$$\left((\lambda)^{(5)}\right)^2+\left((b_{28}')^{(5)}+(b_{29}')^{(5)}-(r_{28})^{(5)}+(r_{29})^{(5)}\right)(\lambda)^{(5)}\Big)$$

$$+\left(\left((\lambda)^{(5)}\right)^2+\left((a_{28}')^{(5)}+(a_{29}')^{(5)}+(p_{28})^{(5)}+(p_{29})^{(5)}\right)(\lambda)^{(5)}\right)(q_{30})^{(5)}G_{30}$$

$$+\left((\lambda)^{(5)}+(a_{28}')^{(5)}+(p_{28})^{(5)}\right)\left((a_{30})^{(5)}(q_{29})^{(5)}G_{29}^{*}+(a_{29})^{(5)}(a_{30})^{(5)}(q_{28})^{(5)}G_{28}^{*}\right)$$

$$\left((\lambda)^{(5)}+(b_{28}')^{(5)}-(r_{28})^{(5)}\right)s_{(29),(30)}T_{29}^{*}+(b_{29})^{(5)}s_{(28),(30)}T_{28}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(6)}+(b_{34}')^{(6)}-(r_{34})^{(6)}\{\left((\lambda)^{(6)}+(a_{34}')^{(6)}+(p_{34})^{(6)}\right)$$

$$\left[\left((\lambda)^{(6)}+(a_{32}')^{(6)}+(p_{32})^{(6)}\right)(q_{33})^{(6)}G_{33}^{*}+(a_{33})^{(6)}(q_{32})^{(6)}G_{32}^{*}\right)\Big]$$

$$\left((\lambda)^{(6)}+(b_{32}')^{(6)}-(r_{32})^{(6)}\right)s_{(33),(33)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(33)}T_{33}^{*}\Big)$$

$$+\left((\lambda)^{(6)}+(a_{33}')^{(6)}+(p_{33})^{(6)}\right)(q_{32})^{(6)}G_{32}^{*}+(a_{32})^{(6)}(q_{33})^{(6)}G_{33}^{*}\Big)$$

$$\left((\lambda)^{(6)}+(b_{32}')^{(6)}-(r_{32})^{(6)}\right)s_{(33),(32)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(32)}T_{32}^{*}\Big)$$

$$\left((\lambda)^{(6)}\right)^2+\left((a_{32}')^{(6)}+(a_{33}')^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\Big)$$

$$\left((\lambda)^{(6)}\right)^2+\left((b_{32}')^{(6)}+(b_{33}')^{(6)}-(r_{32})^{(6)}+(r_{33})^{(6)}\right)(\lambda)^{(6)}\Big)$$

$$+\left(\left((\lambda)^{(6)}\right)^2+\left((a_{32}')^{(6)}+(a_{33}')^{(6)}+(p_{32})^{(6)}+(p_{33})^{(6)}\right)(\lambda)^{(6)}\right)(q_{34})^{(6)}G_{34}$$

$$+\left((\lambda)^{(6)}+(a_{32}')^{(6)}+(p_{32})^{(6)}\right)\left((a_{34})^{(6)}(q_{33})^{(6)}G_{33}^{*}+(a_{33})^{(6)}(a_{34})^{(6)}(q_{32})^{(6)}G_{32}^{*}\right)$$

$$\left((\lambda)^{(6)}+(b_{32}')^{(6)}-(r_{32})^{(6)}\right)s_{(33),(34)}T_{33}^{*}+(b_{33})^{(6)}s_{(32),(34)}T_{32}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(7)}+(b_{38}')^{(7)}-(r_{38})^{(7)}\{\left((\lambda)^{(7)}+(a_{38}')^{(7)}+(p_{38})^{(7)}\right)$$

$$\left[\left((\lambda)^{(7)}+(a_{36}')^{(7)}+(p_{36})^{(7)}\right)(q_{37})^{(7)}G_{37}^{*}+(a_{37})^{(7)}(q_{36})^{(7)}G_{36}^{*}\right)\Big]$$

$$\left((\lambda)^{(7)}+(b_{36}')^{(7)}-(r_{36})^{(7)}\right)s_{(37),(37)}T_{37}^{*}+(b_{37})^{(7)}s_{(36),(37)}T_{37}^{*}\Big)$$

$$+\left((\lambda)^{(7)}+(a_{37}')^{(7)}+(p_{37})^{(7)}\right)(q_{36})^{(7)}G_{36}^{*}+(a_{36})^{(7)}(q_{37})^{(7)}G_{37}^{*}\Big)$$

$$\left((\lambda)^{(7)}+(b_{36}')^{(7)}-(r_{36})^{(7)}\right)s_{(37),(36)}T_{37}^{*}+(b_{37})^{(7)}s_{(36),(36)}T_{36}^{*}\Big)$$

$$\left((\lambda)^{(7)}\right)^2+\left((a_{36}')^{(7)}+(a_{37}')^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\Big)$$

$$\left((\lambda)^{(7)}\right)^2+\left((b_{36}')^{(7)}+(b_{37}')^{(7)}-(r_{36})^{(7)}+(r_{37})^{(7)}\right)(\lambda)^{(7)}\Big)$$

$$+\left(\left((\lambda)^{(7)}\right)^2+\left((a_{36}')^{(7)}+(a_{37}')^{(7)}+(p_{36})^{(7)}+(p_{37})^{(7)}\right)(\lambda)^{(7)}\right)(q_{38})^{(7)}G_{38}$$

$$+\left((\lambda)^{(7)}+(a_{36}')^{(7)}+(p_{36})^{(7)}\right)\left((a_{38})^{(7)}(q_{37})^{(7)}G_{37}^{*}+(a_{37})^{(7)}(a_{38})^{(7)}(q_{36})^{(7)}G_{36}^{*}\right)$$

$$\left((\lambda)^{(7)}+(b_{36}')^{(7)}-(r_{36})^{(7)}\right)s_{(37),(38)}T_{37}^{*}+(b_{37})^{(7)}s_{(36),(38)}T_{36}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(8)}+(b_{42}')^{(8)}-(r_{42})^{(8)}\}\{\left((\lambda)^{(8)}+(a_{42}')^{(8)}+(p_{42})^{(8)}\right)$$

$$\left[\left(\left((\lambda)^{(8)}+(a_{40}')^{(8)}+(p_{40})^{(8)}\right)(q_{41})^{(8)}G_{41}^{*}+(a_{41})^{(8)}(q_{40})^{(8)}G_{40}^{*}\right)\right]$$

$$\left((\lambda)^{(8)}+(b_{40}')^{(8)}-(r_{40})^{(8)}\right)s_{(41),(41)}T_{41}^{*}+(b_{41})^{(8)}s_{(40),(41)}T_{41}^{*}\Big)$$

$$+\left(\left((\lambda)^{(8)}+(a_{41}')^{(8)}+(p_{41})^{(8)}\right)(q_{40})^{(8)}G_{40}^{*}+(a_{40})^{(8)}(q_{41})^{(8)}G_{41}^{*}\right)$$

$$\left((\lambda)^{(8)}+(b_{40}')^{(8)}-(r_{40})^{(8)}\right)s_{(41),(40)}T_{41}^{*}+(b_{41})^{(8)}s_{(40),(40)}T_{40}^{*}\Big)$$

$$\left((\lambda)^{(8)}\right)^2+\left((a_{40}')^{(8)}+(a_{41}')^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\Big)$$

$$\left((\lambda)^{(8)}\right)^2+\left((b_{40}')^{(8)}+(b_{41}')^{(8)}-(r_{40})^{(8)}+(r_{41})^{(8)}\right)(\lambda)^{(8)}\Big)$$

$$+\left(\left((\lambda)^{(8)}\right)^2+\left((a_{40}')^{(8)}+(a_{41}')^{(8)}+(p_{40})^{(8)}+(p_{41})^{(8)}\right)(\lambda)^{(8)}\right)(q_{42})^{(8)}G_{42}$$

$$+\left((\lambda)^{(8)}+(a_{40}')^{(8)}+(p_{40})^{(8)}\right)\left((a_{42})^{(8)}(q_{41})^{(8)}G_{41}^{*}+(a_{41})^{(8)}(a_{42})^{(8)}(q_{40})^{(8)}G_{40}^{*}\right)$$

$$\left((\lambda)^{(8)}+(b_{40}')^{(8)}-(r_{40})^{(8)}\right)s_{(41),(42)}T_{41}^{*}+(b_{41})^{(8)}s_{(40),(42)}T_{40}^{*}\Big)\}=0$$

$$+$$

$$(\lambda)^{(9)}+(b_{46}')^{(9)}-(r_{46})^{(9)}\}\{\left((\lambda)^{(9)}+(a_{46}')^{(9)}+(p_{46})^{(9)}\right)$$

$$\left[\left(\left((\lambda)^{(9)}+(a_{44}')^{(9)}+(p_{44})^{(9)}\right)(q_{45})^{(9)}G_{45}^{*}+(a_{45})^{(9)}(q_{44})^{(9)}G_{44}^{*}\right)\right]$$

$$\left((\lambda)^{(9)}+(b_{44}')^{(9)}-(r_{44})^{(9)}\right)s_{(45),(45)}T_{45}^{*}+(b_{45})^{(9)}s_{(44),(45)}T_{45}^{*}\Big)$$

$$+\left(\left((\lambda)^{(9)}+(a_{45}')^{(9)}+(p_{45})^{(9)}\right)(q_{44})^{(9)}G_{44}^{*}+(a_{44})^{(9)}(q_{45})^{(9)}G_{45}^{*}\right)$$

$$\left((\lambda)^{(9)}+(b_{44}')^{(9)}-(r_{44})^{(9)}\right)s_{(45),(44)}T_{45}^{*}+(b_{45})^{(9)}s_{(44),(44)}T_{44}^{*}\Big)$$

$$\left((\lambda)^{(9)}\right)^2+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\Big)$$

$$\left((\lambda)^{(9)}\right)^2+\left((b'_{44})^{(9)}+(b'_{45})^{(9)}-(r_{44})^{(9)}+(r_{45})^{(9)}\right)(\lambda)^{(9)}\Big)$$

$$+\left((\lambda)^{(9)}\right)^2+\left((a'_{44})^{(9)}+(a'_{45})^{(9)}+(p_{44})^{(9)}+(p_{45})^{(9)}\right)(\lambda)^{(9)}\Big)(q_{46})^{(9)}G_{46}$$

$$+\left((\lambda)^{(9)}+(a'_{44})^{(9)}+(p_{44})^{(9)}\right)\left((a_{46})^{(9)}(q_{45})^{(9)}G^*_{45}+(a_{45})^{(9)}(a_{46})^{(9)}(q_{44})^{(9)}G^*_{44}\right)$$

$$\left((\lambda)^{(9)}+(b'_{44})^{(9)}-(r_{44})^{(9)}\right)s_{(45),(46)}T^*_{45}+(b_{45})^{(9)}s_{(44),(46)}T^*_{44}\Big)\}=0$$

**And as one sees, all the coefficients are positive. It follows that all the roots have negative real part, and this proves the theorem.**

Note: **Same dovetailing explanation holds good for d/dt (partial differential with respect to t), d/dt,d^2/dt^2 (acceralation: double dot). Such an exposition is helpful in optimisation problems, duality of motion, wave motion a homogeneous functions studies, and control theory which we intend to incorporate in future.** It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## Persistence, Coexistence And Thrusty Permanence For Interacting Structured Particles, Transactions In Finance, Econometric Productivity Functions, Environment, Interfaces, Cosmological Interactions Et Al

## Josef Hofbauer And Sebastian J. Schreiber Model Thereof: A Functional Appraisal

ABSTRACT. **JOSEF HOFBAUER AND SEBASTIAN J. SCHREIBER** studied the dynamics of interacting structured systems (in their case population dynamics) can be modeled by $\frac{dxi}{dt}= A_i (x) xi$ where $x_i \epsilon R^n_{i,} x = (x_1,\ldots,x_k)$, and $A_i(x)$ are matrices with non-negative off- diagonal entries. These models are permanent if there exists

a positive global attractor and are robustly permanent if they remain permanent following perturbations of $A_i\ (x)$. Necessary and sufficient conditions for robust permanence are derived using dominant Lyapunov exponents $X_i\ (\mu)$ of the $A_i\ (x)$ with respect to invariant measures $\mu$. The necessary conditions requires $\max_i\ X_i(\mu)>0$ for all ergodic measures with support in the boundary of the non-negative cone. The sufficient condition requires that the boundary admits a Morse decomposition such that $\max_i$ Xi$(\mu)>0$ for all invariant measures $\mu$ supported by a component of the Morse decomposition. When the Morse components are Axiom A, uniquely ergodic, or support all but one population, the necessary and sufficient conditions are equivalent. Applications to spatial ecology, epidemiology, and gene networks are given. Appeared in *Journal of Differential Equations*, 248, 1955-1971 (2010) It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## INTRODUCTION

A fundamental issue in population biology is what are the minimal conditions to ensure the long-term survivorship for interacting populations whether they be viral particles, bio-chemicals, plants, or animals. When these conditions are met the interacting populations are said to persist or coexist. Since the pioneering work of Lotka and Volterra on competitive and predator prey interactions, Thompson, Nicholson, and Bailey on host-parasite interactions, and Kermack and McKendrick on disease outbreaks, nonlinear difference and differential equations have been used to understand conditions for population persistence [1-5]. One particularly important form of persistence for deterministic models is *permanence* or *uniform persistence* which corresponds to the existence of a global attractor bounded away from extinction of one or more species. When such an attractor exists, interacting populations are able to recover from "vigorous shake ups" of the of the population state. Permanence has been characterized from a topological perspective [6-8] and with average Lyapunov functions [9-12] for differential equation models of the form

$$\frac{dxi}{dt} = x_i f_i(x) \qquad i = 1,\ldots,k \qquad (8)$$

Where $x_i$ is the abundance of population $i$, $x = (x_1,\ldots,x_k)$, and $f_i(x)$ is the per-capita growth of population $i$.

An important extension of the concept of the permanence is *robust permanence* [13-16] which requires that (1) remain permanent after sufficiently small perturbations of the per-capita growth rates $f_i$. The importance of robust permanence as one crude form of structural stability for population models are approximations to reality. Consequently, if nearby models (e.g. more realistic models) are not permanent despite the focal model being permanent, and then one can draw few conclusions about the persistence of the biological system being approximated by the model. One can view robust permanence as one crude form of structural stability for debits-credits particles, propositions, transactions in the celestial subterranean realm and ceratoid dualism, finance like debiting say a current account and crediting to another account, or recrediting to the account which had been earlier debited, further conduction of the cash credit account, black hole formation, guzzling of inventories for the output, creation of income from output, reinvestment of income to investment,gormadizing of matter by black holes, and emission of radiation therefore, gamma ray bursts, in fact most of the transactions and interactions in the universe are conservative and quantum mechanical universe, something like a account acting as predator and prey. More often than not as we have seen from Haimovici model such models could be applied to the problems, propositions, which are not subsumed by the concomitant problems in the bearing. the population models. For this perspective, it is not unexpected that there are permanent systems that cannot be approximated by robustly permanent systems [17]. In [14, 16], criteria for robust permanence were developed with respect to the average per-capita growth rates $\int f_i(x)d\mu(x)$ with respect to invariant probability measures $\mu$ supported on the boundary of the positive orthant (i.e. where one or more Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates are extinct). Roughly, these criteria for robust permanence require that the average per-capita growth rate (i.e. invasion rate) is positive for some missing species.

While (1) can account for many types of Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates. Interactions, it assumes that all individual within populations are exactly the same. However, theoretical biologists have long recognized that different individuals within a population may be in different states (e.g. different sizes or ages, living in different spatial locations) and these differences can have important consequences for population dynamics [18-24]. To account for how these differences influence persistence of interacting populations, we develop criteria for robust permanence for the dynamics of $k$ interacting structured populations. Roughly, these dynamics correspond to

replacing $x_i$ in (1) by vectors and $f_i(x)$ by matrices. In section 2, we describe these models in greater detail and review some basic concepts from dynamical systems theory. In section 3, we introduce the structured analog of average per-capita growth rates and define the concept of robustly unsaturated invariant sets. Necessary and sufficient conditions for an invariant set to be robustly unsaturated are proven. For Axiom A or uniquely ergodic invariant sets, the necessary and sufficient criteria are shown to be equivalent. In section 4, we use Morse decompositions and the criteria for robustly unsaturated invariant sets to develop necessary and sufficient conditions for robust permanence. In sections 5 and 6, we provide several applications of our results to spatially structured ecological models, structured epidemiological models, and models of gene networks. We apply the model to econometrics. To give an example investment needs inventory to be used towards the consummation of output, part of output is to be sold for consolidation of income, and some income might be replenished in the investment and so on. In finance a cash credit account could be debited or credited when there is say purchase of demand bills or discount of usance bills. Cash credit account is debited whenever cheques are issued towards the purchase of raw materials and so on. In fact some of the models have been applied to Organic psychiatry and we have shown how such models are useful in dealing with the propositions of Husserl, Derrida, Deleuze and others.

## 8. MODELS AND ASSUMPTIONS

Let $x_i$ denote the state of the $i$-th population that lies in the non-negative cone $C_i$ of $R^{ni.}$ Define $x=(x_1,...x_k)\epsilon C$ to be the non-negative cone of $R^n$ where $n=\sum_{i=1}^{k} n_i$. Let $C_i^+$ denote positive cone { $x_i \epsilon\ C_i$: $\prod_i$i noitalupop rof$^j$

$i$ *and* $C^+=\prod_i C^+_i$ , the positive cone for the interacting populations. If $x \mapsto A_i\ (x)$ is a map into $n_i x n_i$ matrices that describes the growth of population I, then the dynamics of the interacting populations are given by

$$(2)\frac{dx_i}{dt} = A_i(x)x_i \ i=1,...,k$$

Let $x.t$ denote the solution of (2), with the initial condition x, and more generally the semi flow generated by (2). Before stating our assumption on (2), recall a few definitions from dynamical systems. Given sets $I \subseteq R$ and $K \subseteq C$, let $K.I = \{x.t : t \epsilon I, x \epsilon K\}$. A set $K \subseteq C$ *is invariant* if K.t = K for all t>0, and *forward invariant* if $K.t \subset K$ for all $t > 0$. The *omega limit* set of a set $K \subseteq C$ equals $\omega(K) = \bigcap_{t\geq 0,} \overline{K.[t,\infty)}$. The *alpha limit* set of $K \subseteq C$ is $a(K) = \bigcap_{t\leq 0}$

$\overline{K.(-\infty,t]}$. Given a forward invariant set K, $B \subset K$ is an attractor for the semi flow *x.t restricted* to K provided there exists an open neighborhood $U \subseteq K$ of B such that $\omega(U) = B$. The stable set of a *compact invariant* set *K* is defined by

$$W^s(K) = \{x \in C : \omega(x) \neq \emptyset \text{ and } \omega(x) \subseteq K\}.$$

The semi flow generated by (2) is *dissipative* if there exists an attractor B with $W^s(B) = C$.

Throughout this paper, we make the following assumptions:

**A1**: $x \mapsto A_i(x)$ is continuous

**A2:** $A_i(x)$ is irreducible and has non-negative off- diagonal entries,

**A3:** *x.t* is defined for all $t \geq 0$

**A4:** *x.t* is dissipative with global attractor $\Gamma(A)$.

A1 is a basic regularity assumption. The non-negative of off- diagonal entries in A2 implies that there are no negative feedbacks between individuals of different states in population *i*. This assumption is meet for many types of structured models as discussed in sections 5 and 6. The irreducibility assumption A2 is generally meet and implies that all individuals within a population can pass through all states. In Remark1, we discuss how this irreducibility assumption can be relaxed. A3 ensures that the population dynamics are defined for all future time. A4 requires that population densities/ abundances eventually are uniformly bounded, a condition that should be meet for an biologically realistic model.

We say that the semi flow of (2) is *permanent* if this semi flow is dissipative and there is a positive attractor $B \subset C^+$ such that $W^s(B) \subset C^+$. We will denote this positive attractor by $\Lambda(A)$. Permanence ensures that populations can recover from rare large perturbations and allows for a diversity of dynamical behaviors.

For K $\subset$ C and $\delta > 0$, define the $\delta$- *neighborhood* of K as

$$N\delta(K) = \{x \in C : \backslash x - y < \delta \text{ for some } y \in K\}.$$

We define a $\delta$- *perturbation of (2)* to be a system of the form

$$\frac{dx}{dt} = \tilde{A}(x)x$$

that satisfies assumptions A1-A4, $\Gamma(\tilde{A}) \subset N$ *( $\Gamma(A)$), and* $||A(x) - \tilde{A}(x)|| \leq \delta$ *for all* $x \in N$ *($\Gamma(A)$).* (2) is *robustly permanent* if there exist >0 and $\epsilon > 0$ such that all $\delta$- perturbations of (2) are permanent, and $d(\mathcal{A}(\tilde{A}), C \backslash C^{+}) > \epsilon$ for all $\delta$- perturbations i.e., there is a common/ uniform region of repulsion around the boundary.

## INUNDATION OF COMPACT SETS

We begin by studying the linear skew product flows on $\Gamma(A) \times R^{n}_{i}$ defined by $(x,y).t = (x.t, B_i(t,x)y)$ where $Y(t) = B_i(t,x)$ is the solution to $\frac{dY}{dt}(t) = A_i(x.t)Y(t)$ with y(0) equal to the identify matrix. Our assumption that $A_i$ is irreducible with non- negative off diagonal entries implies that $B_i$ (t,x)$C_i \subset C_i^{+}$ for all $x$ and t > 0 (see, e.g., [25] ). A result of Ruelle [26, Prop, 3.2] implies that there exist continuous maps $u_i\, v_i : \Gamma(A) \rightarrow C_i^{+}$ with $u_i(x)| = |v_i(x)| = 1$ such that The line bundle $E_i(x)$ spanned by $u_i(x)$ is invariant i.e. $E_i(x.t) = B_i(x,t)E_i(x)$ for all t ≥ 0.The vector bundle $F_i(x)$ perpendicular to $v_i(x)$ is invariant i.e. $F_i(x.t) = B_i(x,t)F_i(x)$ for all $t \geq 0$.There exist constants $a > 0$ *and* $\beta > 0$ such that(3) $||B_i(t,x)F_i(x)|| \leq a\, exp(-\beta t)||B_i(t,x)E_i(x)||$ For all $x \in K$ *and* $t \geq 0$.**Interaction rates.** Given any x $\epsilon$ $\Gamma$(A), define the *interaction and inundation rate* of *species i at particle/proposition/problem/ state x as* $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right) In$ $||B_i(t,x)||$ when $x_i$ =0, $\lambda_i(x)$ provides an upper bound to the rate of growth of population $i$ when introduced at infinitesimally small densities. Important properties of this invasion rate are summarized in the following two propositions. For instance, the second proposition implies if $(x.t)_i$ stays bounded away from zero (i.e. population $i$ persist), then

$\lambda_i(x) = 0$. Alternatively, if $\lambda_i(x) < 0$,then population $i$ is doomed to extinction. **Proposition 1.** $\lambda_i(x)$ *satisfies the following properties:* $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right) In \ |B_i\ (t,x)v|$ *for any* $v > 0$ *in* $R^{ni}$, *and* $\lambda_i(x) = limsup_{t\to\infty}\left(\frac{1}{t}\right) \int^{t}_{0} (A_i(x.s)u_i(x.s), u_i(x.s))ds.$

**Proof:** To prove the first property, we first show that

$$\lambda i(x) = \underset{t\to\infty}{limsup} \frac{1}{t} In\ |B_i(t,x)u_i(x)|.$$

To this end, let $v \in R^{ni}$ be any non-zero vector. Since $R^{ni} = E_{i(x)}$ $F_i(x)$, we can write $v = au_i(x) + w$ with $a \in R$ and $w \in F_i(x)$. Equation (3) implies

$|B_i(t,x)v| \leq a|B_i(t,x)u_i(x)| + |B_i(t,x)w| \leq |B_i(t,x)u_i(x)|(a + \alpha + exp(-\beta t)|w|)$. Hence, $\limsup_{t\to\infty} \frac{1}{t} In\, B_i(t,x)\, v| \leq$

$$\limsup_{t\to\infty} \frac{1}{t} In\, B_i(t,x)\, u_i(x) \leq \lambda_i(x)$$

for all non zero vectors $v \in R^{ni}$. Since $||B_i(t,x)|| = sup|v|=1 B_i(t,x)v|$, this inequality implies that $\lambda i(x) = \limsup_{t\to\infty} \frac{1}{t}$ In $|B_i(t,x)v|$, this inequality implies that $\lambda_i(x) = limsup_{t\to\infty} \left(\frac{1}{t}\right)$ In $|t,x)u_i(x)|$. Now let v $\epsilon$ $C_i^+$. Then, we can write $v = au_i(x) + w$ with $a > 0$ and $w \in F_i$.

Equation (3) implies

$$|B_i(t,x)v| \geq a|B_i(t,x)u_i(x)| - |B_i(t,x)w|$$

$$\geq |B_i(t,x)u_i(x)|(a - \alpha + exp(-\beta t)|w|).$$

Since a>0, this inequality implies that

$$\lambda i(x) \geq \limsup_{t\to\infty} \frac{1}{t} \text{In } |B_i(t,x)v| \geq \lambda i(x) \geq \limsup_{t\to\infty} \frac{1}{t} \text{In } |B_i(t,x)u_i(x)| = \lambda i(x)$$

To prove the second assertion, let $b_i(t,x) = In\, |B_i(t,x)u_i(x)|$. Invariance of $u_i(x)$ implies that $b_i(t,x)$ is additive:

$$B_i(t+s,x) = In|B_i(t+s,x)\, u_i(x)| = In\, |B_i(s,x.t)B_i(t,x)u_i(x)$$

$$= \text{In } |B_i(s,x.t)\, u_i(x.t)||\, B_i(t,x)\, u_i(x)|$$

$$= \text{In } b_i(s,t.x)\, b_i(t,x)$$

Additivity of $b_i(t,x)$ and the fact that $b_i(0,x) = 0$ implies

$$\frac{d}{dt} b_i(t,x) = \lim_{s\to 0} \frac{bi(t+s,x - bi(t,x)}{S-}$$

$$= \lim_{s\to 0} \frac{bi(s,x.t)}{S} = \frac{d}{ds}|_{s=0}\, b_i(s,x.t)$$

$$=$$

$$= (A_i(x.t)u_i(x.s), u_i(x.s),$$

The Fundamental Theorem of Calculus implies

$$B_i(t,x)= \int_0^t (A_i(x.s)\, u_i(x.s), u_i(x.s))ds\backslash$$

and the second assertion follows.

**Proposition 2:** *For the solutions of (2) we have:* if $x_i>0$, then $\lambda i(x) \leq 0$. If $\lambda i(x)<0$, then $\lim_{t\to\infty} (x.t)_i=0$ If $limsup_{t\to\infty}$ $(x.t)_i|>0$, then $\lambda i(x)=0$.

**Proof:** First, assume that x satisfies $x_i>0$. Since the semi flow is dissipative, there exists $\gamma >0$ such that $|x.t| \leq \gamma$ for all $t \geq 0$. Proposition 1 and the definition of the skew product flow imply

$$\lambda i(x) = \limsup_{t\to\infty} \frac{1}{t} \text{In } |B_i(t,x)x_i| = \limsup_{t\to\infty} \frac{1}{t} \text{In } | (x.t)_i |$$

$$\leq \limsup_{t\to\infty} \frac{In\,\gamma}{t} = 0$$

To prove the second assertion, assume that $\lambda i(x) < 0$. $If\ x_i=0$, then the invariance of the faces of C imply that $(x.t)_i=0$ for all $t\geq 0$. Alternatively, if $x_i> 0$, then Proposition I and the skew product flow imply

$$\limsup_{t\to\infty} \frac{1}{t} In|(x.t)_i| = \lambda i(x) < 0.$$

The final assertion follows from the first two assertions. 5

**3.2. Invariant measures.** We review some definitions from ergodic theory. Given a Borel probability measure $\mu$ on ***C***, the *support* of $\mu$, denoted supp( $\mu$),is the smallest closed set whose complement has measure zero. A Borel probability measure $\mu$ is called *invariant* for (2) provided that Integral $h(x)d\mu(x) =$ Integral $h(x.t)d\,\mu(x)$for all $t \geq 0$ and for all bounded continuous functions $h$:C→R. An invariant measure $\mu$ is called *ergodic* provided that $\mu(B)$=0 or 1 for any invariant Borel set *B.*

For an invariant measure $\mu$ define the invasion rate species $i$ with respect to $\mu$ as

$$\lambda_i(\mu) = \int_c \langle A_i(x)u_i(x), u_i(x)\rangle d\mu(x)$$

**Proposition 3:** Let $\mu$ be an *invariant measure.* $\lambda_i(\mu)$ *Satisfies the following properties:* $\lambda_i(x) = \lim_{t\to\infty} \frac{1}{t}\int_0^t \langle A_i(x.s)u_i(x.s), u_i(x.s)\rangle, ds$ exists for $\mu$ - *almost every* $x$. *Moreover, if* $\mu$ *is ergodic,* $\lambda_i(x) = \lambda_i(\mu)$ for $\mu$-*almost every* $x$.*If* $\mu$ *is ergodic, then there exists* $I \subset \{1, \dots, k\}$such that $\mu(\Pi_{i\in I} C_i^+)$=1 and $\lambda_i(\mu) = 0$ *for all* $i \in I$.

**Proof:** Let $\mu$ be an invariant measure. Define $h_i(x) = \langle A_i(x)u_i(x), u_i(x)\rangle$. The Birkhoff's' ergodic theorem and Proposition 1 imply

$$\lambda_i(x) = \lim_{t\to\infty} \frac{1}{t}\int_0^t h_i(x.s)ds$$

Exists for $\mu$- almost every $x$. Moreover, $\lambda_i(x) = \lambda_i(\mu)$ $\mu$- almost surely if $\mu$ is ergodic.

Assume $\mu$ is ergodic. By ergodicity and invariance of the faces of **C**, there exists $I \subset \{1, \dots, k\}$ such that $\mu(\Pi_{i\in I} C_i^+)$=1. Let $K \subset \Pi_{i\in I} C_i^+$ be a compact set such that $\mu(K) > 0$. The Poincare recurrence theorem and the Birkhoff ergodic theorem imply that there is $x \in \Pi_{i\in I} C_i^+$such that $x.t_n \in K$ for some $t_n \uparrow \infty$ and $\lambda_i(x) = \lambda_i(\mu)$ for all $i$. The third assertion of Proposition 2 implies $\lambda_i(\mu) = 0$ for all $i \in I$.

**3.3. Robustly unsaturated sets**. Let $K \subset \partial C$ be a compact isolated invariant set for the flow $x.t$ restricted to $\partial C$.*K* is unsaturated if $W^s(K) \subset \partial C$ and K is isolated for $x.t$. If *K* is not unsaturated, then *K* is saturated. *K* is robustly unsaturated for (2) if there exists $\delta > 0$ such that the continuation of K for any $\delta -$ perturbation of (2) is unsaturated.**Theorem 1:** Let *K* be a compact isolated invariant set for $x.t$ restricted to$\partial C$. If one of the following equivalent conditions hold For all invariant measures $\mu$ supported by *K* (4) $\max_{1\le i\le k} \lambda_i(\mu) = 0$

There exist $p_i > 0$ such that (5) $\sum_{1\le i\le k} p_i \lambda_i(\mu) = 0$

For all ergodic probability measures supported by *K*.

then *K* is robustly unsaturated for (2). Alternatively, if $x \to A(x)$ is twice continuously differentiable and *K* is robustly unsaturated, then (4) holds for all ergodic measures $\mu$ supported by *K*.

**Remark 1**: For some applications (e.g., the disease model considered in section 6.1) it useful to relax the irreducibility assumption A2. For instance, if there exists an open neighborhood U of K such that for each $i$, $A_i(x)$ has a fixed off-diagonal sign pattern for all $x \in U$,then $A_i(x)$ can be decomposed into a finite number, say $m_i$, of

irreducible components. For each of these irreducible components, one can define $\lambda_i^j(\mu) = \int_C \langle A_i^j(x)u_i^j(x), u_i^j(x)\rangle \, d\mu(x)$ where $A_i^j(x)$ is the submatrix of $A_i(x)$ corresponding to the $j$-th irreducible component of $A_i(x)$ and $u_i^j(x)$ is the continuous invariant sub bundle for the irreducible linear cocycle determined by $A_i^j(x)$. If we define $\lambda_i(\mu) = \max_{1\le j\le m_i} \lambda_i^j(\mu)$, then all of the assertions of Theorem 1 still hold.

**Proof:** To prove the first assertion of the theorem, we need the following lemma. We call an invariant measure $\mu$ for (2) saturated if $\lambda_i(\mu) \le 0$ for all $i$

**Lemma 1**: Let $\delta_n$ be a non –negative sequence that converges to zero as $n\to\infty$. If $\mu_n$ are saturated invariant measures for $\delta_n -$ perturbations of (2), then the weak* limit points of $\{\mu_n\}_{n=1}^{\infty}$ is a non-empty set consisting of saturated invariant measures for (2).

**Proof:** Let $\delta_n$ be a non-negative sequence that coverages to zero. Let $\mu_n$ be saturated invariant measures for $\delta_n$ perturbations, $\dot{x} = A^n(x)x$, of $\dot{x} = A(x)x$. By weak* compactness of Borel probability measures supported on $N_1(\Gamma(A))$, there exist weak* limit points of $\mu_n$. Let $\mu$ be such a weak* limit point. Since $\mu_n$ are supported by $N_1(\Gamma(A))$, for all $n$ sufficiently large, $\mu$ is supported by $N_1(\Gamma(A))$.To verify that $\mu$ is an invariant measure for (2),let $h: C \to R$ be a bounded continuous function. Let $t > 0$ and $\in > 0$ be given. Choose $n$ sufficiently large so that $|\int_c h(x)\, d\mu(x) - \int_c h(x) d\mu_n(x)| \le \in$, $|\int_c h(x.t)\, d\mu(x) - \int_c h(x.t) d\mu_n(x)| \le \in$, and $|h(x.t) - h(x^n.t)| \le \in$ for all $x \in N_1(\Gamma(A))$.Then $|\int_c h(x.t) - h(x) d\mu(x)|$ is

$\le \quad |\int_c h(x.t)\, d\mu(x) - \int_c h(x.t)\, d\mu_n(x)| + |\int_c h(x.t)\, d\mu_n(x) - \int_c h(x)\, d\mu(x)|$

$\le \quad \in + \int_c |h(x.t) - h(x^n.t)|\, d\mu_n(x)| + |\int_c h(x^n.t)|\, d\mu_n(x) - \int_c h(x) d\mu(x)|$

$\le \quad 2 \in + |\int_c h(x^n.t) - h(x)\, d\mu_n(x)| + |\int_c h(x) d\mu_n(x) - \int_c h(x) d\mu(x)|$

$\le \quad 3 \in$

Where the last line follows from the invariance of $\mu_n$ for $x^n.t$. Since $\in > 0$ and t > 0 are arbitrary, we have $\int_c h(x.t)d\mu(x) = \int_c h(x)d\mu(x)$ for all t >0 and all bounded continuous functions $h: C \to R$. It follows that $\mu$ is an invariant measure for (2). To see that $\mu$ is saturated, define $h_i^n(x) = \langle A_i^n(x)u_i^n(x), u_i^n(x)\rangle$ where $u_i^n(x)$ spans the invariant one-dimensional bundle given by $[26, Prop.3.2]$ for $\frac{dx}{dt} = A^n(x)x$. Since $A_i^n(x) \to A_i(x)$ uniformly for $x \in N_1((A))$ as $n \to \infty$, $[26, Prop.3.2]$ implies that $u_i^n(x) \to u_i(x)$ converges uniformly for $x \in N_1(\Lambda(A))$. Given $\in > 0$, choose $n$ sufficiently large so that $|h_i^n(x) - h_i(x)| \leq \in$ for all $x \in N_1(\Lambda(A))$ and $|\int_c h_i^n(x)\,d\mu(x) - \int_c h_i^n(x)\,d\mu_n(x)| \leq \in$.

Then

$$\lambda_i(\mu) \leq \int_c |h_i(x) - h_i^n(x)|d\mu(x) + |\int_c h_i^n(x)\,d\mu(x) - \int_c h_i^n(x)\,d\mu_n(x)|$$

$$+ \int_c h_i^n(x)\,d\mu_n(x)$$

$$\leq 2\in$$

Since $\in > 0$ is arbitrary, $\lambda_i(\mu) \leq 0$ for all $i$ and $\mu$ is saturated.

Using this lemma, we prove that if K is saturated, then there exists a saturated invariant measure $\mu$ supported by K. Assume K is saturated. Work of Hofbaue and So $[7, Thm.2.1]$ implies that either there exists $y \in W^s(K) \cap C^+$ or K is not isolated for the unrestricted flow $x.t$ If there exists $y \in W^s(K) \cap C^+$, then for all $t \leq 0$ define $v_t = \frac{1}{t}\int_0^t \delta_{y.s}\,ds$ where $\delta_{y.s}$ denotes a Dirac measure based at the point $y.s$. Dissipativeness of (2) and weak* compactness of the Borel probability measures supported by (A) imply there exist $t_k \to \infty$ such that the sequence $v_{tk}$ converges in the weak* topology to a Borel probability measure $\mu$ with support in K.A standard argument implies that $\mu$ is $x.t$ invariant. Define $h_i: C \to R$ by $h_i(x) = (A_i(x)u_i(x), u_i(x))$.Proposition 2 and weak* convergence imply that

$$0 \geq (\lambda i(y) = \limsup_{t\to\infty} \frac{I}{t}\int_0^t h_i(y.s)ds$$

$$= \limsup_{t\to\infty} \int_c h_i\,dv_t$$

$$\geq \limsup_{t\to\infty} \int_p h_i\,d\,v_{tk} = \lambda i(\mu)$$

For all $i$. Alternatively suppose that K is not isolated for the semi flow Then there exists a sequence of positive $\omega - limit$ sets that accumulate on K. Let $\mu_n$ be a sequence of ergodic probability measures supported by these $\omega -$ limit sets. Propositions 1 and 3 imply that $\mu_n$ are saturated for all $n$ Applying Lemma I with $\delta_n$=0 for all $n$ implies that there exists a saturated invariant measure $\mu$ suppoorted by *K*. Hence we have shown that if *K* is saturated then there exists a saturated invariant probability measure $\mu$ supported by *K* Equivalently, (4) holding for all invariant probability measures with support in *K* implies that *K* is unsaturated.

Next we show that if *K* not robustly saturated, then there exists a saturated invariant measure supported by *K* Indeed suppose *K* is not robustly saturated, Then there exists a non-negative sequence $\delta_n$ converging to zero and a sequence of saturated measures $\mu_n$ for $\delta_n$-perturbations of (2)with support in the continuation of *K*. Let $\mu$ be a weak$^*$ limit point of $\{\mu_n\}\infty_{n=1}$ Lemma 1 implies that $\mu$ is saturated. Moreover, since the continuation of *K* converges to *K* as $\delta_n \to 0, \mu$ is a saturated invariant measure for (2) supported by *K* Hence we have shown that if *K* is not robustly saturated then there exists a saturated invariant measure supported by *K* Equivalently if (4) holds for all invariant measures with support in *K* then *K* is robustly unsaturated.

To see the equivalence of the conditions given by (4) and (5) let $\Delta = \{p \in R_+^k : \sum_i P_i = 1\}$ and notice that

$$\text{Min max } \lambda i\ (\mu = \min_{\mu} \max_{p \epsilon \Delta} \sum_i Pi\lambda i(\mu)$$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K* The Minimax theorem (see, e. g.,[27] implies that

$$(6)\ \min \max \lambda i(\mu) = \min_{\mu} \max_{p \epsilon \Delta} \sum_i Pi\lambda i(\mu)$$

Where the minimum is taken over invariant probability measures $\mu$ with support in *K*. Since min$\mu$

$\sum_i P_i \lambda_i(\mu)$ is attained at an ergodic probability measure with support in *K*, the equivalence of the conditions given by (4) and (5) is established.

To prove the final assertion of the Theorem assume $x \to A(x)$ is twice continuously differentiable and there exists a saturated ergodic measure $\mu$ supported by *K* Proposition 3 implies that there exists *I*C {1,…,k} such that (□$i \in iC_i^+$)=1and $\lambda i(\mu) = 0$ For all $i \in I. SinceKc\,.\partial$□,$\{1, \ldots . k\}\backslash I$ is non-empty. We will show that (2) is not robustly permanent by proving that for all $\delta > 0$ there exists a $\delta$ −perturbation of (2) for which the continuation of *K* is saturated Let $\delta > 0$ be given Choose V $c$ N$\delta$ (*Γ*(A)) to be a compact neighborhood of Γ(A) such that V.t $c$ int *V* for all t > 0.Let $\eta > 0$ be such that *W*=:*N* $\eta$(Γ(A))$c$ *V* Let p:c→ $[0,1] be$ a$C^{\infty}$function such that $p(x) = 1 \quad for\ all\ x \in \Gamma(A) and\ p(x) = 0 for\ all\ x \in C\backslash W\ Define\ \tilde{A} = (\tilde{A}_1, \ldots, \tilde{A}_k)$ by

$$\tilde{A}_i(x) = \{ \frac{A(x)}{A(x) - \frac{\delta}{2Idp\,(x)}} \quad \frac{if\ i \epsilon I}{if\ i \notin I}$$

Where Id denotes the identity matrix of appropriate dimension Let x. t denote the semi flow of $x = \tilde{A}(x)x$.Since$x.t = x.t$ $whenever\ x.[0,t] \in C\backslash W, it\ follows\ that \quad \tilde{V}.t \subset intV\ for\ all\ t >$ $0. Hence, \Gamma(\tilde{A})c\ V\ c\ N\delta\ (\ (A))). We\ also\ have\ ||A(x) - \tilde{A}(x)|| \leq \frac{\delta}{2}$Therefore $x = \tilde{A}(x) is\ a\ \delta -$ perturbation of (2) By construcation,$x.t = x.t\ for\ all\ x \in K and\ t \geq 0. Hence, \mu$ is ergodic for the semiflow of $\dot{x} = \tilde{A}(x)x\ and\ \lambda_i(\mu) \leq -\frac{\delta}{2}$ for this semiflow and $i \in \{1, \ldots, k\}\backslash I$.Let $\mathcal{L}$ and O be the Lyapunov exponents and Oseledec regular points supported by $\mu$ for $x = \tilde{A}(x)x(see, e.g.,$

[28] for definitions).At each point $x \in O$, the splitting of $R^n$ determines three subspaces: the stable subspace $E^s(x)$, the center subspace $E^c(x)$, and the unstable subspace $E^u(x)$ Proposition 3 and our choice of Ã imply that $E^s(x) \cap C^+ \neq \emptyset$. The Pesin stable manifold theorem [28, Corollaries 3.17 and 3.18] implies that tangent to $E^s(x)$ $E^c(x)$ and $E^u(x)$ are locally $x.t$ −invariant families of $C^1$ discs $W^s_x, W^c_x$ and $W^u_x$ corresponding to the stable ,center and unstable manifolds The family of stable manifolds $W^s_x$ is contained in $W^s$ (Ã,supp $(\mu)$). $Since\ E^s(x) \cap C^+ \neq \emptyset$, $W^s_x \cap C^+ \neq \emptyset$ for some $x \in K$. Consequently *K* is saturated for $x = \tilde{A}(x)x$.

The difference between the sufficient and necessary condition is that for the sufficient condition (4) has to hold for all invariant measures supported by *K*, while for the necessary condition (4) has to hold only for ergodic measures supported by *K*. since the invariant measures lie in the convex hull of the ergodic measures the sufficient condition can be more restrictive then the necessary condition.

For three Classes of invariant sets the necessary and sufficient conditions coincide Recall a compact invariant set *K* for (2) is uniquely ergodic if K only supports one invariant measure Recall a compact invariant set *K* is Axiom A if the flow of (2) restricted to *K* is transitive and hyperbolic (see, e.g.,[29] for definitions.)

**Corollary 1:** Assume $x \to A(x)$ is twice continuously differentiable It K$\subset \partial$□ is a compact invariant set and either is uniquely ergodic or Axiom A, then *K* is robustly unsaturated if and only if (4) holds for all ergodic, measures $\mu$ with support in *K*

Proof. If K is uniquely ergodic then the assertion follows immediately For an Axiom A invariant set Sigmund [30, thm.1] has proven that invariant measures supported by periodic orbits of *K* are dense in the set of invariant probability measures supported by *K* In particular any invariant measure can be approximated by an ergodic measure and consequently the result follows.

Another special case where the necessary and sufficient conditions coincide is when *K* supports all populations except one for discrete time models this case was considered by [22].

**Corollary2:** Assume $x \to A(x)$is twice continuously differentiable Let K$c$ $II_i \geq 2C_i^+$ be a compact invariant set. Then K is robustly unsaturated if and only if $\lambda_1(\mu) > 0$ for all invariant measures supported by *K*

## 4. MORES DECOMPOSITIONS AND ROBUST PERMANENCE

To state the sufficient condition for robust permanence, we use a characterization of permanence due to Garay [8] and Hofbauer and So [7] that involves Morse decompositions of the boundary flow Conley[31] defined a collection of sets$\{M_{1,\ldots,}M_m\}$ to be a Morse decomposition for a compact invariant set *K* if $M_1\ldots,M_m$ are pairwise disjoint, compact isolated invariant sent for the flow of (2) restricted to K. For each $x \in K$ there are integers r=r$(x) \leq s =$

$s(x)$ such that $a(x) \subseteq M_r$and $\omega(x) \subseteq M_s$. If $r(x) = s(x), then\ x \in M_r$. Garay, Hofbauer and so [7,8] proved the following characterization of permamnce

**Theorem 2** (Garay, Hofbauer-so) if $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$then (2) is permanent if and only if each of the components $M_1$ are unsaturated.

Theorem 1 and 2 imply the following result: As stated in the above all these results are true for most of problems, interactions, transactions in cosmology, econometrics, finance and other subjects, which we shall illustrate as we go on constructing models for the same.

**Theorem 3:** If $\{M_1,\ldots..,M_m\}$ is a Morse decomposition for $\partial$and (4) holds for each of the components of the More decomposition, then (2) is robustly permanent Conversely if $x \to A(x) is\ twice\ continuously\ differentiable$

*and (4) is violated by an ergodic measure supported by one of the components of the Morse decomposition, then (2) is not robustly permanent.* Theorem 3 in conjunction with Corollaries 1 and 2 yield a characterization of robust permanence for a class of structured models.

**Corollary 3:** *Let $\{M_1,\ldots..,M_m\}$ os a Morse decomposition for* $\partial$C. *Assume $x \mapsto A(x)$ is twise continuously differentiable and for each Morse component $M_i$ one of the following assertions hold $M_i$ is Axiom A, $M_i$ is uniquely ergodic, or There exists $j \in \{1,\ldots,k\}$ such that $M_i$ ----- $\{x \in C : x_l \gg 0 \text{ for } l \neq j\}$ Then (2) is robustly permanent if and only if (4) holds for all ergodic measures supported by $U_i M_i$.*

Our result also provided a structured analogue for characterizing totally permanent systems[32]. For an ergodic probability measure$\mu$, define sup ($\mu$) to be the subset $I$ -----$\{1,\ldots,k\}$ such that $\mu(---_{i \in I} C_+) = 1$

**Corollary 4:** *The following statements are equivalent: For all ergodic probability measures $\mu$ with support in* $\partial$C, *$\lambda_i(\mu) > 0$ for all $I \in \{1,\ldots,k\} \setminus supp(\mu)$ (2) and all of its subsystems are robustly permanent.*

*Proof. Suppose the first statement holds. Then unsaturated condition given by (5) with p= (1,…,1) and Theorem 3 implies that (2) and all its subsystems are robustly permanent.. The other direction follows form the necessary condition for robust permanence in Theorem 3 and the second assertion of Proposition 3.*

*To illustrate the broad applicability of theorem 3, we develop applications to spatially structured models in section 5 and to a disease model and a gene network model in section 6.*

## 5. APPLICATIONS: PATCH MODELS

A fundamental application of our result are to spatially structured models, with *k* species dispersing between *m* patches {33-35}:

(7) $\frac{dx_i^j}{dt} = x_i^j f_i^j(x^j) + \sum d_i^{jl} x_i^l - e_i^j x_i^j$ i = 1,..,k j=1,…,m

*l*

Where $x_i^j$ denotes the density of species I in patch *j*, $x^j = (x^j_1,.....,x^j_k)$ is the vector of species densities in patch *j*, $f^j_i$ is the per-capita growth rate of species *i* in patch *j*, $d_i^{jl} \geq 0$ is the dispersal rate for species *i* from patch *l* into patch *j*, and $e_i^j$ is the emigration rate of species *i* out of patch *j*. Hence $e_i^j \geq \sum_l d_i^{lj}$. We assume that the matrices $(d_i^{lj})_{j,l}$ are irreducible for each species *i* . it is then easy to write (7) in the form (2), with $x_i = (x^l_i, ..., x_i^m)$ 'where' denotes transpose.

For a single species (7) generates a monotone flow. Under mild assumption, e.g., each $f^j_i$ *is* decreasing and negative for large densities, there is a globally stable equilibrium, see, e.g.,[36] and [37, sect. 5.4]. This equilibrium will be positive if the invasion rate at the origin $\lambda(0)$, which is given by the leading eigenvalue (stability modulus) of the matrix $(d^{jl} + f^j(0) - e^j)$, is positive.

5.1 Two types of debitable and creditable systems:. For *two competing species*, e.g., $f^j_i$ is decreasing with respect to both species and is negative at large densities, the dynamics are still monotone, and hence almost all orbits converge to an equilibrium. Robust permanence requires two conditions. First, $\lambda_i(0) > 0$ for both species in which case there are two single species equilibria $E_1$ and $E_2$. Second, the invasion rates $\lambda_2(E_1)$ and $\lambda_2(E_2)$ are positive. Whether there is a unique positive globally stable equilibrium for the robustly permanent system depends in a delicate way on the

system parameters, see [38]. For slowly dispersing populations (i.e. $0< e^j_i \ll 1$ for all *i, j* ), the robust permanence condition is particularly

Straight forward to verify. Let $E^j_i$ be the largest solution to $x^j_i \quad f_i{}^j \ (x^j_i) = 0$ i.e. the equilibrium attained by species *I* in patch *j* when there is no dispersal and no competitors. For slowly dispersing populations, $E_1$ is close to the product $(E_1{}^1, \ldots, E_1{}^m)$, and the matrix $A^2(E_1)$ is close to the diagonal matrix $\text{diag}(f^i{}_2(E_1{}^j))^m_{j=1.}$ Hence the invasion rate $\lambda_2(E_1)$ is close to $\max_j f_2^j(E^j_1)>0$. So we obtain: The system with two slowly dispersing competing species is robustly permanent if and only if system refers to system of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

(8) $\max_j f_2^j(E^j_1)>0$ and $\max_j f_2^j(E^j_2)>0$

Intuitively, robust permanence requires for each species there is at least one patch where it can persist.

For *systems, particle interactions, financial transactions, cosmological interactions and reactions, as we shall show model by model the* situation are even easier. Since (under mild assumption on the prey dynamics) the Morse decomposition of $\partial C$ is given by two equilibria $\{0,E_1\}$, the spatial predator-prey system is robustly permanent, if the predator can invade at the prey equilibrium: $\lambda_2(E_1) > 0$. However, the global dynamic are likely to be more complicated. For instance, if in each patch there is a globally stable limit circle (e.g.,[39-41]) and dispersal is sufficiently slow (i.e. $0< e^j{}_i \ll 1$for all *i, j*), our result about robust permanence and the theory of normally hyperbolic manifolds [42] imply that there is a positive *m* dimensional torus, which attracts almost all orbits in C, whenever $\max_j f_2^j(E^j_1)>0$

**5.2 Rock-Paper-Scissors Dynamics.** The Lotka-Volterra model of rock-paper-scissor dynamics is a simple model that is used as property for understanding intransitive ecological outcomes [43,44]. Here, a simple spatial version of this dynamic is given by

$$\frac{dx_1^j}{dt} = x_1^j(1 - x_1^j - \beta^j x_2^j - \alpha^j x_3^j) + \sum d_1^{jk} x_1^k - e_1^j x_1^j$$

$\frac{dx_2^j}{dt} = x_2^j(1-\alpha^j x_1^j - x_2^j - \beta^j x_3^j) + \sum d_2^{jk} x_2^k - e_2^j x_2^j$

$\frac{dx_3^j}{dt} = x_i^j(1-\beta^j x_1^j - \alpha^j x_2^j - \beta^j x_3^j) + \sum d_3^{jk} x_3^k - e_3^j x_3^j$

Where $a^j \in (0,1)$ and $\beta^j > 1$ for all $j$. A more general version of this dynamic is presented in [44]. Under the assumption that there is no cost to dispersal (i.e. $e^j_i = \sum_l d^{l,j}_i$ for all $i, j$) , the maximal invariant set in $\partial C$ consist of the origin 0 and a heteroclinic cycle connecting positive single species equilibria $E_1$, $E_2$ and $E_3$. For slowly dispersing population ($0 < e^j_i \ll 1$ for all $i$, j), $E_1$ is close to $(x_1, x_2, x_3) = (1,\ldots,1,0,\ldots,0,0,\ldots,0)$ and the invasion rates $\lambda_2(E_1)$ and $\lambda_3(E_1)$ are close to $\max_j 1-a^j$ and $\max_j 1-\beta^j$. For the other equilibria $E_i$, the debitable rates of the missing species are also given by $\max_j 1-a^j$ and $\max_j 1-\beta^j$. Consider the Morse decomposition of $\partial C$ given by 0 and the heteroclinic cycle. Since $\lambda i(0) = 1 > 0$ for $i = 1, 2, 3, 0$ is robustly insaturated. An algebraic computation reveals that (4) holds for all invariant supported by the heteroclinic cycle (i.e. all convex combination of the Dirac measures supported by the equilibria) if and only if

$$\max_j(1-a^j) > -\max_j(1-\beta^j)$$

Equivalently,

$$2 > \min_j a^j + \min_j \beta^j$$

In particular, even if the heteroclinic cycle is attracting for each patch when the system is uncoupled (i.e. $2 < a^j + \beta^j$ for all $j$), it can be repelling for the weakly coupled system.

5.3. Three species **Lotka Volterra in spatially homogenous environments.** Consider Lotka-Volterra dynamics in a spatially homogenous environment: $(x+a)^n =$

(9) $$\frac{dx_i^j}{dt} = x_i^j\left(r_i - \sum_s a_{is}x_s^j\right) + \sum_l d_i^{jl} x_i^l - e_i^j x_i^j, \qquad i = 1,\ldots,k;\ j = 1,\ldots,m$$

We assume that dissipative can be shown by a linear Liapunov function: There are $c_i > 0$ such that $\sum_i c_i\, a_{is} x_i x_s < 0$ holds for all $x \in R_+^k \setminus \{0\}$. Then the weighted sum of densities across all patches $\sum_{i,j} c_i\, x_i^j$ is decreasing for large densitities and hence (9) generates a dissipative semiflow with global attractor $r$. We further assume that for each patch *j* and each species *i* the immigration rate equals the emigration rate , i. e,

(10) $$e_i^j = \sum_l d_i^{jl}$$

This guarantees that the set $H = \{x \in C: x_i^j = x_i^l \forall i, j, l\}$ of spatially homogenous states is forward invariant under (9).

Now consider the Lotka-Volterra dynamics without spatial structure

(11) $$\frac{dx_i}{dt} = x_i(r_i - \sum_s a_{is} x_s), \qquad i = 1, \ldots, k$$

(11) $$\frac{dx_i}{dt} = x_i(r_i - a_{is} x_s), \quad i = 1, \ldots, k$$

It has been shown [12,ch.16.1,16.2] That for $k = 3$,(11) is robustly permanent (robust meaning here within the class of Lotka –Volterra systems) if and only if all equilibria on the boundary of $R_+^3$ are unsaturated and whenever there is a heteroclinic cycle connecting the one species equilibria then this cycle is repelling. Equivalently the following four conditions hold. There exists an interior equilibrium $\dot{x}$ (i.e $A\dot{x} = r \;\; with\; \dot{x} \gg 0$); Det (-A)>0; the 2 species subsystems are not bistable competition systems if there is a heteroclinic cycle between the one species equilibria, say $\mathbf{E_1} \to \mathbf{E_2} \to \mathbf{E_3} \to \mathbf{E_1}$ then the following inequality holds

12) $$\lambda_2(E_1)\lambda_3(E_2)\lambda_1(E_3) > |\lambda_3(E_1)\lambda_1(E_2)\lambda_2(E_3)|$$

Furthermore, the proof in [12, ch.16.1, 16.2] shows that the boundary flow of (11) has a simple Morse decomposition: if (iv) applies then there are two Morse sets, the origin 0 is a repeller, and the heteroclinic cycle is

the dual attractor (within$\partial R^3_+$); if (iv) does not apply then the finitely many boundary equilibria form a Morse decomposition of the boundary flow.

**Species, population we mean types of and population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.**

**Theorem 4:** Under the above assumptions,(9) for k=3 is permanent if and only if *(i),(ii),(iii)* and *(iv)* hold.

Proof: Suppose (9) is robustly permanent. Since the nonspatial system (11) is the restriction of (9) to the invariant subspace **H**, (11) must be robustly permanent and hence*(i)-(iv)* hold

Conversely, suppose that (i)-(iv) hold. We show that the above Morse decomposition of the boundary flow of (11) (now in **H**) is also Morse decomposition for (9) restricted to the maximal invariant subset of $\partial C$ which is contained in the spatial two species subsystems of (9). For this we show that in each 2 species system the orbits of (9) converge to **H** (For the one species subsystems this is obvious.) If a two species subsystem of (11) has an internal equilibrium, say $\mathbf{E_{12}}$ then by assumption (ii),this is not a saddle but a sink, and hence $a_{11}a_{22} > a_{12}a_{21}$. This implies that the 1-2 submatrix of A is VL-stable [12,ch.15.3,],and [45] implies that $\mathbf{E_{12}}$ is the global attractor for the spatially structured system (9) restricted to the first two species.

If a two species subsystem of (11) has no internal equilibrium, then a one species equilibrium say $\mathbf{E_1}$ is the global attractor. Then Lemma 2 (see below) show the corresponding result for (9).A similar proof applies to the case where the origin **0** is the global attractor of a two species subsystem.

Finally observe that under the above assumptions (9)-(10), the invasion rates $\lambda_i$ at spatially homogenous boundary equilibria are the same for (9) and (11) Since all boundary equilibria are unsaturated for (11), they are for (9) as well. The same applies in case (iv) to the heteroclinic cycle. Note that the condition (12) on the eigenvalues is equivalent to the fact that all invariant measures supported on the heteroclinic cycle are unsaturated, see [16,Ex.4.5]

Species we mean types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

**Lemma 2:** Suppose, in a two species Lotka –Volterra system (11) (k=2) all interior orbits converge to the one species equilibrium $\mathbf{E_1}$ .Then all orbits of the spatial version (9) (k=2) converge to the spatially homogeneous one species equibrium $\mathbf{E_1} \in \mathbf{H}$.

**Proof:** If the system is of predator-prey or mutualistic type then Hasting's [45] result applies again. So let us assume that the local interaction is given by the competition system

(13) $\frac{dx_1}{dt} = r_1 x_1(1 - x_1 - ax_2)$ $\frac{dx_2}{dt} = r_2 x_2(1 - \beta x_1 - x_2)$

With 0<a <1< $beta$.In this case, Kazuo Kishimoto (letter to JH, Oct 1987) has given the following argument. There exists a family of forward invariant rectangles contracting to $\mathbf{E_1}$. This shows that all interior solutions of (9) (k=2) converge to $\mathbf{E_1}$.

Theorem 4 shows that a spatial network of identical patches is permanent, if the within patch dynamics is permanent, and three species Lotka-Volterra dynamics. It is not clear how to extend this result to more than three species. The crucial step in the three species case is that permanence precludes bistable two species subsystems. The spatial version of a bistable two species system allows plenty of stable equilibria outside **H,** see [33].However, permanent four species Lotka Volterra systems may have two species bistable subsystems [12,ch.16.4].So there may be invariant sets on the boundary outside **H** which need to be unsaturated to make (9) permanent.

## 6. MORE APPLICATIONS

6.1. **Disease, particle, product, marketing, financial, econometric dynamics with density-dependent demography**. GAO and Hethcote [46] introduced a model of disease dynamics with density-dependent demography. Here we consider a variation of their model in which a population of size P has a per-capita birth rate b(P) and per –capita death rate d. To ensure that the population persists in the absence of the disease, we assume that there exists K>0 such that b(K)=d. To describe the disease dynamics, let *S* be the number of individuals susceptible to the disease, *I* the Number of infected individuals, and R the number of individuals that have recovered from the disease. We assume that P=S+*I* +R. On the same lines we can think of sick and non performance accounts, and those which are conducting themselves well, and those which are satisfactory and find a parallel formula. The same thing holds true in case of a black hole eating matter and black hole radiating radiation.. If the disease transmission is asymptotic, β is the contact rate between susceptible and infected individuals, γ is the rate at which individuals recover from the disease, and m the mortality rate due to disease, then the population dynamics is given by

$$\frac{dp}{dt} = b(P)P - dP - mI$$

$$\frac{dI}{dt} = \beta(P - I - R)\frac{I}{\epsilon + P} - (\gamma + d + m)I$$

$$\frac{dR}{dt} = \gamma I - dR$$

Where $\epsilon>0$ is a constant It is useful to introduce a change of coordinates in which $y=\frac{I}{P}$ and $z=\frac{R}{P}$:

$$\frac{dP}{dt} = (b(P)-d-my)P$$

$$\frac{dy}{dt} = \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(P)\ y$$

$$\frac{dz}{dt} = \gamma y - (b(P) - my)z$$

Setting $x_1$=P and $x_2$=(y,z) 'where' denotes transpose yields a structured model where $A_1(x)$=b-d-my and

$$A_2(x) = \begin{pmatrix} \beta(1-y-z)\frac{P}{\epsilon+P} - \gamma - m(1-y) - b(p) & 0 \\ \gamma & my-b(p) \end{pmatrix}$$

Since {0,(*K*,0,0)}is a Morse decomposition of the boundary dynamics Theorem 3 and Remark 1 imply that this model is robustly permanent if and only if b(0)>d and $\beta\frac{K}{\epsilon+K} > \gamma + m + d$. In particular, for $\epsilon > 0$ sufficiently small, one requires that the basic reproductive number $\frac{\beta}{r+m+d}$ is greater then one.

Remarkably this same criterion determines robust permanence of model with significantly more complicated boundary dynamics. For example, suppose the focal population is a predator species. If the prey has abundance *N* and exhibits logistic dynamics f(*N*)is the per-capita predator consumption rate of the prey and b(N)is the per-capita reproductive rate of the predator then the dynamics become

$$\frac{dN}{dt} = rN(1 - N/K - f(N)P$$

$$\frac{dP}{dt} = (b(N) - d - my)P$$

$$\frac{dy}{dt} = \beta(1 - y - z)\frac{P}{\epsilon+P} - r - m(1-y) - b(N)\ y$$

$$\frac{dz}{dt} = \gamma y - (b(N) - my)z$$

Where r>0 and *K*>0, and *f*(*N*),b(N)satisfy f(0)=b(0)Define $x_1 = N$, $x_2 = p$, $x_3 = (y,z)'$ Let us assume that the disease-free community (i.e.y=z=0)is permanent (i.e.b(K)>d) and let A be the global attractor in the interior of the P-N plane Under these assumption a Morse decomposition of the boundary dynamics is given by $(x_1, x_2, x_3) = (0,0,0)$, $(x_1, x_2, x_3) = (K, 0,0)$and A. Since the equilibria (0,0,0) and (k,0,0) are unsaturated, it remains to characterize $\curlywedge_3(\mu)$for any invariant measure μ supported by A. Proposition 3 and Remark 1 imply

$$\curlywedge_2(\mu) = \int_c (bx_1)\text{-}d)d\mu(x) = 0$$

*Hence* $= \int_c bx_1)d\mu(x) = d$ *and*

$$\curlywedge_3(\mu) = \beta \int \frac{x_2}{\epsilon + x_2} d\mu(x) - r - m - d$$

**Theorem 3**: and Remark 1 implies that this structured model is robustly permanent if $\frac{\beta}{r+m+d} > 1 and\ \epsilon > 0$ is sufficiently small.

6.2. Gene networks. The repressilator is an oscillatory gene network based on three (or more generally an odd number of transcriptional repressors. Such a system has been genetically engineered in E-coli by Elowitz and Leibler [47] Mathematical model go back to [48.49] they all involve concentrations of proteins and mRNs. A modified model with auto-activation was suggested in [50]

$$(14)\ x_i = \beta(y_i - x_i)$$

$$(15)\ y_i = aF(x_{)i} \text{-} y_i)$$

With

$$(16)\quad F(x)_i = x_i g\ (x_i, x_i - 1) = \frac{x_i}{1 + x_i + p x_{i-1} + K p\, x_i x_i - 1}$$

Here $x_i$ and $y_i$ are normalized concentrations of proteins and mRN As belonging to gene $i$,p is the strength of repression and k a cooperativity parameter this is a structured system of type(2)

$$(17)\ \begin{pmatrix} x_i \\ y_i \end{pmatrix} dot = \begin{pmatrix} -\beta & \beta \\ ag(x_i,x_i)-1 & \end{pmatrix} \begin{pmatrix} x_i \\ y_i \end{pmatrix}$$

In [50,Thm 3] it was shown that for n=3,a>1 and p>1,system(14)has a heteroclinic cycle connecting the 3 single gene equilibria (similar to the rock –scissors-paper dynamics of section5.2)Moreover

- the system is permanent i.e. the boundary of $R^6$+is repelling, if $\lambda + \mu > 0$
- this heteroclinic cycle is asymptotically stable if $\lambda + \mu < 0$

where $\lambda > 0$ and $\mu < 0$ are the invasion rates of the two missing genes i.e. the leading eignvalues of the 2X2 matrix in (17)evaluated at the appropriate single gene equilibria (for explicit expressions see[50,eqs.(127,128)].)

Theorem 3 implies the permanence condition implies the stronger conclusion of robust permanence A similar result holds for n odd, but the heteroclinic cycle is then actually a heteroclinic network.

Note: For references please see the original paper.

One of the examples for which we apply the model are as follows: **Progress in Particle and Nuclear Physics Volume 88, May 2016, Pages 1–28 Chiral magnetic and vortical effects in high-energy nuclear collisions—A status report D.E. Kharzeeva, b, J. Liaoc, d, , , S.A. Voloshine, G. Wangf**

The interplay of quantum anomalies with magnetic field and vorticity results in a variety of novel non-dissipative transport phenomena in systems with chiral fermions, including the quark–gluon plasma. Among them is the Chiral Magnetic Effect (CME)—the generation of electric current along an external magnetic field induced by chirality imbalance. Because the chirality imbalance is related to the global topology of gauge fields, the CME current is topologically protected and hence non-dissipative even in the presence of strong interactions. As a result, the CME and related quantum phenomena affect the hydrodynamical and transport behavior of strongly coupled quark–gluon plasma, and can be studied in relativistic heavy ion collisions where strong magnetic fields are created by the colliding ions. Evidence for the CME and related phenomena has been reported by the STAR Collaboration at Relativistic Heavy Ion Collider at BNL, and by the ALICE Collaboration at the Large Hadron Collider at CERN. The goal of the present review is to provide an elementary introduction into the physics of anomalous chiral effects, to describe the current status of experimental studies in heavy ion physics, and to outline the future work, both in experiment and theory, needed to eliminate the existing uncertainties in the interpretation of the data. Keywords Chiral magnetic effect; Chiral vortical effect; Chiral anomaly; Quark-gluon plasma; Heavy ion collisions

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# Sufficient Conditions For The Existence And Linear Stability Of Time Independent Age Distributions In Respect Of Particles, Propositions, Problems, Hegelian Structures, Deleuzean Treatises, Interactions And Transactions In Cosmos And Others

**Abstract:** Gopalaswamy studied sufficient conditions are obtained for the existence and linear stability of time independent age distributions in two species competition with age time lagged density dependent mortality and fertility functions.. Work was published in Bulletin Australian Mathematical Society Volume 25 (1982). We show how the model could be applied in case of finance, econometrics, psycho economics, interactions between particles and transactions in nature. Basically the competition is for resources, notwithstanding such experiments could be conducted in laboratory like a rat kept in front of a cat, such incidents could be represented mathematically, and the corresponding linear stability conditions ca be obtained in case of interactions in cosmos. Take for example the case of a cash credit account of a borrower. Many inward bills waiting for paid, or inland or foreign bills payment all are in line for the debit of the account, which also acts as a resource. In addition there are recurrent cases where in the credit is issued to the cash credit account in case of demand and inland bills. In case of economics, investment has inventories and it has to be utilized towards the end of consummation of output. And this is case of output gobbling part of investment which is inventories. Output has to be sold to the customer and the amount received is diverted partly to savings and consumption and part could be employed towards the consolidation of the system. So this is a case of investment itself acting as a both receiver and payer. It shall eat and also produce the finance. Savings as stated above could be deployed for the further diversification and proliferation and that is towards the end of putting the company on terra firma. Another example in cosmology is the case of black holes guzzling matter. There are lot of black holes and lot of matter. Result of such a guzzling the production of Hawking radiation. Horava gravity is another example. How the conformal field theory could be changed in the case of Horava gravity. Still another case is the case of Vonneumann algebra in the cae of Horava gravity. In this case like cash credit account, the Horava gravity and Von Neumann algebra act as both receivers and takers. Such incidents are legion in cosmology and even in philosophy. We have considered self as individual general ledger, and the debits and credits are raised to the same account thus anointing as both a receiver and payer. We give plethora of such cases and apply the model. Due to spatial constants, same module is used for all the modules and we rewrite to state that the model is applicable to all modules. This is due to spatial restrictions. What is to be done is that each and every parameter is to be changed with a subscript or a superscript and finally it is to be concatenated towards the end of obtention of prediction model or optimizing model, for Duality of motion model for further analysis. Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a

qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

**Introduction**

Let $P_1$ *(t)* and $P_2$ *(t)* denote the total particles, interactions, transactions, etc, which we dilate upon further investigation. Perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. **Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems** sizes (or biomasses) at time $t \geq 0$ of two interacting species living in a common habitat and competing for a common pool of resources. The competition for resources will be implicit in our model similar to that in the two species Lotka- Volterra competition system. Assuming constant sex ratios in the two species we can consider $P_1$ *(t)* and $P_2$ *(t)*to be the Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

population of f; immigration, emigration and internal dispersion in the habitat are assumed to play no significant role in the dynamics of the community.Popoulation we mean Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

We suppose that the two species contain respectively $p_1(a, t)\,da$ and $p_2(a, t)\,da$ individuals with ages between $a$ and $a + da$ $(a \geq 0)$ at time $t$ so that we have

(1.22) $P_1(t) = \int_0^\infty \rho_1(a, t)\,da,\ P_2(t) = \int_0^\infty \rho_2(a, t)\,da,\ t \geq 0.$

The rates of change of the two species densities are by definition $D\rho_1(a, t)$ and $D\rho_2(a, t)$ where

(1.23) $D\rho_i(a, t) = \limsup_{h \to 0} \frac{\rho_i(a+h, t+h) - \rho_i(a,t)}{h},\ i = 1, 2.$

Assuming the existence of the limits in (1.2) we consider the following time lagged model system

(1.24) $$\begin{cases} D\rho_1(a,t) + f_1(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_1(a, t) = 0, & a>0,\ t>0 \\ D\rho_2(a,t) + f_2(a, P_1(t), (a, P_1(t), P_2(t-\tau))\rho_2(a, t) = 0, \text{'} \\ (1.4)\quad \rho_1(0, t) = \int_0^\infty b_1(a, P_1(t), P_2(t-\tau))\rho_2(a, t)da, & t>0, \\ \rho_2(0, t) = \int_0^\infty b_2(a, P_1(t), P_2(t-\tau))\rho_1(a, t)da, \end{cases}$$

(1.5) $\rho_i(a, s) = \varphi_i(a, s),\ a.0,\ s \in [-\tau, 0],$ $i = 1, 2,$

(1.6) $P_i(s) = \int_0^\infty \varphi_i(a, s)\,da,\ s \in [-\tau, 0],$

Where $\tau$some is fixed nonnegative constant; $f_1$ and $f_2$ denote the age and density dependent mortality rates with time lags in interspecific interactions, $b_1$ and $b_2$ denote the age and density dependent fertility functions again having time lags in the interspecific interactions; $\varphi_1$ and $\varphi_2$ denote the initial age distributions needed for the formulation of the model. Note we mean the population of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

The model system (1.1)-(1.6) will be meaningful only if $\varphi_1$, $\varphi_2$, $f_1$, $f_2$, $b_1$, $b2$ are nonnegative functions of the respective arguments. Also since the total initial populations have o be finite, $\varphi_1(\cdot, s)$ should belong to $L_1(R^+)$ for each $s\epsilon$ $[-\tau, 0]$. That $f_i$ and $b_i$ $(i = 1, 2)$ are independent of t explicitly and this assumption can be interpreted to represent the temporal uniformity of the environment.

One of the fundamental questions for (1.1)-(1.6) is the following; under what conditions on $f_i$, $b_i$ and $\tau$ does the system (1.1)-(1.4) have time independent nontrivial solutions, and when such solutions, and when such solutions exist, are they stable with respect to some suitable stability criterion? When such stable time invariant solutions exist for (1.1)-(1.4) we say that the two types has a stable age distribution. Existence of a stable age distribution is the analogue of the existence of a stable steady state for the corresponding age independent system. Types we mean types of Particles, predications, postulations, ions, quarks, molecules, interactions, transact=ions, production rates, and destruction rates or mortality rates.

In this article we investigate the above question for the case of a competitive interaction with some additional constraints on the vital rates $f_i$ and $b_i$ $(i = 1, 2)$. Age dependent population systems without time delays have been considered by several authors (Gurt in and MacCamy [3, 4, 5], Gurt in and Levine [6], Haimovici [7] and Rotenberg [8]). The question of the relation between an age independent system and a corresponding age dependent system has been considered by Gurt in and MacCamy [4] and Gurt in and Levine [6] who have established an asymptotic (*as* $t\rightarrow\infty$ ) relation between such models by constructing a higher dimensional 1 umped parameter system (in terms of ordinary differential equations) to represent the age dependent distributed parameter system. Assuming the existence of stationary age distributions Haimovici [7] considers their stability in a system of two interacting populations explicitly taking into consideration the dynamical nature of the encironments resources and pollution.

In this article we add another realistic feature namely time lags in the age dependent models and obtain sufficient conditions for the existence of stationary age distributions in (1.1)-(1.4) and show that the time lags in interspecific

interactions have no effect on the linear stability of the age distributions although the decay rates of perturbations will depend on the magnitudes of the time delays. Such "harmless" natures of time lags in interspecific interactions have been noted by the author for the age independent models in [1, 2].

## 8. Existence and uniqueness of solutions

We first narrate our assumptions on the vital rates of $f_i$ and $b_i$ $(i = 1, 2)$ so that the system (1.1)-(1.6) will denote a competition system. Let $BC\ (0, \infty)$ denote the linear space of bounded continuous real functions on $(0, \infty)$ with the norm defined by

$$||f|| = \{max\ |\ f(t)\ |,\ t\ \in (0, \infty)\}.$$

Let $K_1$ be the cone of nonnegative functions of $BC\ [0, \infty]$. Let $C\ [-\tau, 0]$ denote the space of real continuous functions with a norm defined by $||g|| = \{\ max\ |g(s)\ |,\ s \in [-\tau, 0]\}$; let $K_2$ be the nonnegative cone of functions in $C\ [-\tau, 0]$ . We present our first set of assumptions on $f_i$, $b_i$ and $\varphi_i$ for $i = 1, 2$.

$A_1$. $f_1 \in C\ (R^+ \times K_1 \times K_2, R^+)$, $f_1 \geq d_1^*$ on $R^+ \times K_1 \times K_2$

$f_2 \in C\ (R^+ \times K_2 \times K_1, R^+)$, $f_2 \geq d_2^*$ on $R^+ \times K_2 \times K_1$ ,

( $d_1^*$ , $d_2^*$ are positive constants);

$f_1$ and $f_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ ; that is, *for* $i = 1, 2$,

$$|f_i\ (a, P_1^{'}\ , P_2^{'}) - f_i(a, P_1^{"}\ , P_2^{"})| \leq d_i\ [||P_1^{'} - P_1^{"}, P_2^{'} - P_2^{"}||]$$

($d_1$, $d_2$ are positive constants),

$$\partial f_1/\ \partial P_1, \partial f_1/\ \partial P_2 \in C(R^+ \times K_1 \times K_2, R^+)\ ,$$

$$\partial f_2/\ \partial P_1, \partial f_2/\ \partial P_2 \in C(R^+ \times K_2 \times K_1, R^+),$$

$A_2$. $b_1 \in C\ (R^+ K_1 \times\ K_2, R^+)$, $b_1 \leq \beta_1^* on\ R^+ \times K_1 \times K_2$ ,

$b_2 \in C\ (R^+ K_2 \times\ K_1, R^+)$, $b_2 \leq \beta_2^* on\ R^+ \times K_2 \times K_1$ ,

*($\beta_1^*$ , $\beta_2^*$ are positive constants);*

$b_1$ , $b_2$ are Lipschitzian with respect to their last two arguments uniformly in $a \in R^+$ *; that is for* $i = 1, 2$,

$$|b_i\ (a, P_1^{'}, P_2^{'}) - b_i\ (a, P_1^{"}, P_2^{"})| \leq \beta_i\ [||P_1^{'} - P_1^{"}|| + ||P_2^{'} - P_2^{"}]$$

$B_1$ , $B_2$ Are positive constants),

$b_i(\circ,(P_1,P_2)\epsilon L^{\infty}(R^+,R^+)$ For fixed $P_1,P_2,i=1,2,$

$$\frac{\partial b_1}{\partial P_2},\frac{\partial b_1}{\partial P_2}\epsilon\ C\ R^{+}\text{x}\ k_{\ 1}\text{x}\ k_{\ 2},\text{R}^{-}),$$

$$\frac{\partial b_2}{\partial P_1},\frac{\partial b_2}{\partial P_2}\epsilon\ C\ R^{+}\text{x}\ k_{\ 2}\text{x}\ k_{\ 1},\text{R}^{-}),$$

$A_3.\ \varphi_i\ (\cdot\ ,\text{s})\in \text{L}^{\ 2}\ (\text{R}^{+},\text{R}^{+})\ \cap \text{C}(\text{R}^{+},\text{R}^{+})$ for each $s\in\ [-\tau,0],$

$i,\ =1,\ 2,$

By a solution of $(1.1)-(1.6)$ on $[0,\text{T}]$, we mean

$\text{P}_i\text{: R}^{+}\ \square[-\tau,\text{T}]\rightarrow \text{R}^{+}\ \ (i,\ =1,\ 2,)$ With the following properties:

$S_1.\quad \text{P}_i(\cdot\ ,t)\in \text{L}^{\ 1}\ (\text{R}^{+},\text{R}^{+})$ for each $t\ \in[-\tau,\text{T}];$

$S_2.\quad \text{P}_i(a,\cdot\ )\in \text{C}\ ([-\tau,\text{T}],\text{R}^{+})$ for each $a\ \in \text{R}^{+}$;

$S_3.\quad \text{DP}_i$ Exists along the characteristics $t-a=$ constant

On $\text{R}^{+}\ \square\{\text{R}^{+}\cap\ [0,\text{T}]\}$ And is continuous $\ (i,\ =1,\ 2,);$

$S_4.\quad \text{P}_i$ Satisfies the system (1.1)- (1.6) for

$(a,t,)\ \in \text{R}^{+}\ \square[-\tau,\text{T}]$ .

Let us first convert the system (1.1)- (1.6) into an equivalent system of integral equations. For brevity we denote by P the pair ( $\text{P}_1,\text{P}_2$) If we let

$$\text{P}_i(a+x,t+x)=\quad \text{P}_i(x),(i,\ =1,\ 2,)\ ,$$

Then (1.3) considered at $(1.3)$ considered at $(a+x,\ t+x)$ becomes

$$\frac{dp_{\ i}(\text{x})}{dx}+f_i(a+x,\ P(t+x))\ \ P_i(x)=0$$

Which has a solution of the form?

$$P_i(x)=P_i(0)\ \exp\left[-\int_0^{\infty} f_{\ i}\big(a+s,P(t+S)\big)ds\right]$$

And hence

$$(2.1)\ p_{\ 1}\ (a+x,t+x)=p_{\ 1}\ (a,t)\ exp\big[-\int_o^{x}\ f_{\ 1}\ \big(a+s,p(t+s)\big)ds\big]$$

For $a \geq t$ *we choose (a,t)=(a-t,0)* in (1.3) and *x=t* in (2.1) so that

(2.2) $p_1(a,t) = \varphi_1(a-t,0)\exp\left[-\int_o^t f_1\left(a-t+s, p_1(s)\right)ds\right]$

$$O \leq t \leq a, i = 1,2,$$

With

$$P_1(s)=(p_1(s),p_2(s-t)),$$

$$P_2(s)=(p_1(s-t),p_2(s)),$$

For $0 \leq a < t$ we choose *(a,t)=(0,t-a)in (1.3)and x=a in (2.1)* so that

(2.3) $p_i(a,t)= B_i(t-a)exp[-\int_o^a f_1(s, p_i(t-a+s))ds]$

$$t > a \geq, i = 1,2,$$

Where

(2.4) $B_i(t) = p_i(0, t), i = 1,2, t > 0.$

Define $M_i and\ L_i\ (i = 1,2)$ as follows:

$$M_i(a,t, p_i(t)) = exp\left[-\int_o^{t-a} f_1(s, p_i(a-a+s))ds\right]$$

(2.5) $L_i\left(a,t,\ p_i(t)\right) = \exp\left[-\int_o^t f_1(\text{a+s},\ p_i(s))ds\right]$

It will now follow from (1.1),(1.3),(2.2)-2.5)that

(2.6) $p_i(t) = \int_o^t B_1(a)\, M_1(a, t, p_i(t))da + \int_o^\infty \varphi_i(a, 0)\ L_i(a,t,p_i(t)da,$

$$t > 0, i = 1,2,$$

(2.7) $p_i(s) = \Phi_i(s) = \int_o^\infty \varphi_i(a, s)da, s \in [-\tau, 0], i = 1,2,$

(2.8) $\beta_i(\text{t}) = \int_0^t \beta_i(\text{a}) b_i(\text{t-a}, P_i(t)) M_i.\left(a,t,P_i(t)\right)da$

$$+\int_0^\infty b_i\left(a+t, P_i(t)\right)\varphi_i(a,0)L_i\left(a,t,P_i(t)\right)da$$

$$i = 1, 2, \text{t} > 0$$

the equivalence of (2.6)-(2.8) with (1.1)-(1.6) is established by the following whose proof is identical to a similar result of Gurtin and MacCamy [4]; hence we omit the details of the proof of the following.

THEOREM 1 : *Suppose the system* (1.1)-(1.6) *satisfies the conditions* $A_1, A_2, A_3$ . *If* $(P_1, P_2)$ *is a solution of*

*(1.1)-(1.6) then the total population* $P_1, P_2$ *and the birth rates* $B_1, B_2$ *satisfy the integral equation (2.6)-(2.8).*

*conversely if* $P_1, P_2, B_1, B_2$ *are nonnegative continuous solution of (2.6)-(2.8) on [ 0, T] and if* $P_1, P_2$ *are*

*defined by (2.2)-(2.4) on* $R^+$ x [0, T] *then such* $P_1, P_2$ *provide a solution of (1.1)-(1.6) on* $R^+$ x [0, T] .

The following a *priori* estimates are useful to prove our existence theorem below. It will immediately follow from the bounds for $f_i$ and $b_i$ that

$$M_i(a,t,P) \leq \exp[-d_i^*(t-a)] \leq 1, if \quad t \geq a,$$

*(2.9)*

$$L_i(a,t,P) \leq \exp[-d_i^* t]$$

From (2.7)-(2.8) we derive that

$$\beta_i(t) \leq \int_o^t \beta_i^* \beta_i(a) \exp[-d_i^*(t-a)]\, da + \beta_i^* \Phi_i(0) \exp[-d_i^* t], \quad i = 1,2,$$

And hence by Gronwall's inequality,

(2.10) $$\beta_i(t) \leq \beta_i^* \Phi_i(0) \exp[\delta_i t], \ \delta_i = \beta_i^* - d_i^*$$

(2.5), (2.6) and (2.10) lead to

(2.11) $$|P_i(t) - 0_i(0)| \leq \beta(0) exp[\delta_i a - d_i^*(t-a)]\, da$$

$$+\int_0^\infty \varphi_i(a,0)|1 - exp(-d_i^* t)|da$$

$$\leq \beta_i^* \Phi_i(0) exp(\beta_i^* t)\left\{\frac{1-exp[-(d_i^*+\beta_i^*)t]}{d_i^*+\beta_i^*}\right\} + \beta(0) d_i^* t$$

$$\leq [\beta_i^* \beta(0) exp(\beta_i^* t) \beta(0) d_i^*] t \quad for \ t > 0.$$

We can now prove the following result on the existence and uniqueness of solution of the integral equations (2.)-(2.8)

**THEOREM 2:** *Assume that the system* (1.1)-(1.6) *satisfies the hypotheses*

| |
|---|
| $A_1, A_2, A_3$. *Then there exists a unique set of continuous function* $(P_1, P_2, B_1, B_2)$ such *that* |
| $P_i : R^+ \to R^+, \qquad B_i : R^+ \to R^+ \ (i = 1,2\ ),$ |
| *Which satisfy the integral equations* (2.6)-(2.8) |
| Proof. We will first show the existence of local solution on same (possibly small) interval [0, T] $\subset R^+$ such that |
| $P_i \in \mathrm{C}([0,T], R^+), \quad B_i \in \mathrm{C}([0,T], R^+)\ , i = 1, 2,$ |
| Satisfying (2.6)-(2.8) and then show the $P_i$ and $B_i$ can be continued as solution for any finite interval in $R^+$ for some fixed |
| Positive number $T$ we let |
| $C^+[-\tau, T] = \{f \in \mathrm{C}([-\tau, T], R)\,|f \geq 0\}$ |
| Let $\|.\|_T$ denote the supremum norm in $\mathrm{C}^+[-\tau, T]$. Then for any |
| $(x, y) \in \mathrm{C}^+[-\tau, T] \times \mathrm{C}^+[-\tau, T]$ we define |
| $\|(x,y)\| = \|x\|_T + \|y\|_T,$ |
| $\mathrm{C}_T = \mathrm{C}^+[-\tau, T], \quad \mathrm{C}_T^2 = \mathrm{C}^+[-\tau, T] \times \mathrm{C}^+[-\tau, T].$ |
| |
| Members of $\mathrm{C}_T$ will be denoted by $Z_t$ with the meaning that |
| $Z(t+s) = z_t(s)$ *for* $s \in [-\tau, T]$ *and* $t \in [0, T]$. Now for any fixed |
| $P_t = (P_{1t}, P_{2t}) \in \mathrm{C}_T^2$, (2.8) is a linear system of uncoupled Volterra integral equations in B = $(B_1, B_2)$ and hence (2.8) will |
| have a unique continuous solution under the assumption $A_1 - A_3$. let us denote such a solution by |
| $\mathrm{B}(t) = B^T(P_t), \quad t \in [0, T]$ |
| Or |
| $B_1(t) = B_1^T(P_{1t}, P_{2t}), \quad B_2(t) = B_1^T(P_{1t}, P_{2t}), \ t \in [0, T]$ |
| If we supply this solution $(B_1, B_2)$ in (2.6) we find that (2.) will be satisfied if and only if $(P_{1t}, P_{2t})$ is a fixed point of the |
| operator |

$$\Pi(P_t) = \{\Pi_1(P_{1t}, P_{2t}),\ \Pi_2(P_{1t}, P_{2t})\}$$

Where $\pi_1$ *and* $\pi_2$ respectively denote the right sides of (2.6) for *i =1,2* .

It is not difficult to see from the nature of $M_i$ , $L_i$ and the hypotheses on $f_i, b_i, \varphi_i$ that

$$B: C_T^2 \to C_T^2 \text{ and } \Pi C_T^2 \to C_T^2$$

(which guarantees the nonnegativity of $P_i$ and $B_i$). We will first show that

There exist positive constants $r^*, t^*$ such that $\pi$ is a mapping of a sphere

$$S_{t*}(\Phi, r^*) = \{ (P_{1t}, P_{2t}) \in C_{t*}^2 \mid \|P_t - \Phi_t\|_{t*} \le r^* \}$$

Into itself and also is a contraction where $\Phi_t$ is defined on [-$\tau$, $t^*$]

By the following:

$$\Phi_t = (\Phi_{1t}, \Phi_{2t});$$

(2.12)

$$\Phi_{it} = \begin{cases} \int_0^\infty \varphi_i(a, s)\, da \text{ for } s \in [-\tau, 0]. \\ \int_0^\infty \varphi_i(a, 0)\, da \text{ for } s \in [0, t^*]. \end{cases}$$

We note that similar to (2.11) one gets the estimates

$$|\Pi(P)(t) - \Phi_i(0)| \le \{B_i^* \Phi_i(0) \exp[B_i^* t] + \Phi_i(0) d_i^*\} t$$

Which together with (2.7) and (2.12) lead to

(2.13) $$\|\Pi(P_t) - \Phi_t\|_{t*} \le \left(\sum_{j=1}^2 \|\Phi_{jt}\|_{t*} [B_j^* \exp(B_j^* t) + d_i^*]\right) t^*$$

Thus if we choose,

$$r^* = T \sum_{j=1}^2 \|\Phi_{jt}\|_{t*} [B_j^* \exp(B_j^* T) + d_i^*]$$

Then $\Pi$ is a mapping of $S_{t*}(\Phi, r^*)$ into itself for all $t^* \le T$. we then will have to show that for small $t^* > 0$ , $\Pi$ is a contraction.

Let $P_t^{(1)}$ *and* $P_t^{(2)}$ be arbitrarily chosen in $S_{t*}(\Phi, r^*)$ and consider $|\Pi(P_t^{(1)}) - \Pi(P_t^{(2)})|$ for $t \le t^*$.

Corresponding to the chosen $P_t^{(1)}$ , $P_t^{(2)}$ we let $B^{(1)}(a) = B(P_T^{(1)})$ and $B^{(2)}(a) = B(P_T^{(2)})$ . Then from

(2.6),

$$|\Phi(P^{(1)}(t)) - \Phi(P^{(2)}(t)| \leq \left| \int_0^t B_1^{(1)}(a)\left\{M_1\left(a,t,P_t^{(1)}\right) - M_1\left(a,t,P_t^{(2)}\right)\right\}da \right|$$

$$+ \left| \int_0^t [B_1^{(1)}(a) - B_2^{(2)}(a)]\, M_1(a,t,P_t^{(2)})da \right|$$

$$+ \left| \int_0^t \varphi_1(a,0)\left[L_1\left(a,t,P_t^{(1)}\right) - L_1\left(a,t,P_t^{(2)}\right)\right]da \right|$$

$$= J_1 + J_2 + J_3 \text{ (say)}$$

using the elementary inequality

$$|e^x - e^y| \leq |x-y| \ for \ x,y \leq 0$$

We derive from (2.5) that

$$\left| M_1(a,t,\ p_t^{(1)}) - M_1(a,t,\ p_t^{(2)}) \right|$$

$$\leq \int_o^{t-a} \left| f_1(s,\ p_t^{(1)}(a+s)) - f_1(s,\ p_t^{(2)}(a+s)) \right| ds$$

$$\leq d_1 \| p_t^{(1)} - p_t^{(2)} \|_{t^*} \text{(using } A_1\text{)}$$

*Now,using (2.10),*

$$\left| J_1 \right| \leq \int_o^a \beta_1^* \Phi_1(0)\, e^{\delta 1^a} \left(\int_o^{t-a} d_1 \| P_t^{(1)} - P_t^{(2)} \|_t ds\right)da$$

$$\leq d_1 \beta_1^* \Phi_1(0)\, t^2 \| P_t^{(1)} - P_t^{(2)} \|_t e^{\delta 1^t/2}$$

In a similar way

$$\left| J_3 \right| \leq \int_o^\infty \varphi(a,o)\left[\int_o^t \left| f_1(a+s, P_t^{(1)}(s) - f_1(a+s, P_t^{(2)}(s)) \right| ds\right]da$$

$$\leq d_1 \Phi_1(0) \| P_t^{(1)} - P_t^{(2)} \|_t$$

Now we have to estimate $\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right|$ in order to estimate $J_2$, from (2.8),

$$B_1^{(1)}(a) - B_1^{(2)}(a)$$

$$= \int_o^a \left[ B_1^{(1)}(\sigma) - B_1^{(2)}(\sigma) b_1 \left( a - \sigma, P_t^{(1)}(a) \right) M_1(\sigma, a, P_t^{(1)}(a) d\sigma + \int_o^a \left[ B_1^{(2)}(\sigma) \right. \right.$$

$$\text{X} \left[ b_1(a - \sigma, P_t^{(1)}(a) M_1(\sigma, a, P_t^{(1)}(a) - b_1 \left( a - \sigma, P_t^{(1)}(a) M_1 \left( \sigma, a, P_t^{(2)}(a) \right) \right] d\sigma \right.$$

$$+ \int_o^\infty \varphi(\sigma, o) \left[ b_1 \left( \sigma + a, P_t^{(1)}(a) \right) - b_1 \left( \sigma + a P_t^{(2)}(a) \right) \right], L_1 \left( \sigma, a, P_t^{(1)}(a) \right) d\sigma$$

$$+ \int_o^\infty \varphi(\sigma, o) b_1 \left( \sigma + a, P_t^{(2)}(a) \right) \left[ L_1 \left( \sigma, a, P_t^{(1)}(a) \right) - L_1 \left( \sigma, a, P_t^{(2)}(a) \right) \right] d\sigma$$

$$= R_1 + R_2 + R_3 + R_4 \text{ (say)}.$$

By the uniform Lipschitz continuity of $b_1$ in $A_2$ we get

$$\left| b_1 \left( a + \sigma, P_t^{(1)}(a) \right) - b_1 \left( \sigma + a, P_t^{(2)}(a) \right) \right| \leq B_1 \, \| P_t^{(1)} - P_t^{(2)} \| \text{ for } a \, \varepsilon [o, t^*]$$

And hence we have

$$|R_3(a) + R_4(a)| \leq \beta_1 \Phi_1(0) || P_t^{(1)} - P_t^{(2)} ||_{t^*} + \Phi_1(0) \beta_1 d_1 || P_t^{(1)} - P_t^{(2)} ||_{t^*}^{*a},$$

$$|R_2(a)| = [\int_0^a [B_1^{(2)}(\sigma) \{ b_1( a-\sigma, P_t^{(1)}(a)) - b_1(a-\sigma, P_t^{(2)}(a)) \} M_1(\sigma, a, P_t^{(1)}(a))$$

$$+ b_1(a, -\sigma, P_t^{(2)}(a)) \{ M_1(\sigma, a, P_t^{(1)}(a)) - M_1(a-\sigma, P_t^{(2)}(a)) \}] d\sigma]$$

$$+ b_1 \left( a - 0, P_t^{(2)}(a) \right) \{ M_1(\sigma, a, P_t^{(1)}(a)) - M_1(a-\sigma, P_t^{(2)}(a)) \}] d\sigma]$$

$$\leq \int_0^a [\beta_1^* \Phi_1(0) \, e^{\delta 1^a} \beta_1 \| P_t^{(1)} - P_t^{(2)} \|_{t^*} + \beta_1 d_1 \| P_t^{(1)} - P_t^{(2)} \|_{t^*} (a-o)] do$$

$$\leq \beta_1^* \Phi_1(0) \, e^{\delta 1^a} \| P_t^{(1)} - P_t^{(2)} \|_{t^*} \{ \beta_1 a + d_1 a^2 / 2 \}$$

Thus

$$\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right|$$

$$\leq \int_o^a \left| B_1^{(1)}(\sigma)\text{-}B_1^{(2)}(\sigma) \right| B_1^{*d}\sigma$$

$$+ \Phi_1(0) /\!/ P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} \left\{ B_1 + B_2\, d_1\, a + B_1^{*} \left( B_2 a + d_1 a^2/2 \right) e^{\delta 1^a} \right\}$$

$$\leq \beta_1^* \int_o^a \left| B_1^{(1)}(\sigma)\text{-}B_1^{(2)}(\sigma) \right| \text{do} + \left( r^* + /\!/ \Phi_t /\!/_{t*} \right) \| P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} g_1^{(a)}$$

Where $g_1(a) \rightarrow \beta_1$ as $a \rightarrow 0$. By Gronwall's inequality it will follow that

$$\left| B_1^{(1)}(a)\text{-}B_1^{(2)}(a) \right| \leq \left( r^* + /\!/ \Phi_t /\!/_{t*} \right) /\!/ P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} \{ g_1(a) + 0(a) \}$$

Using this estimate for $R_1(a)$ we get

$$|R_2(a)| \leq \left| \int_0^a [B_1^{(2)}(\sigma)\, B_1^{(2)}(\sigma),]\, b_1(a)) \text{-} b_1(a\text{-}\sigma, P_t^{(1)}(a))\, M_1(\sigma, a, P_t^{(1)}(a)) do \right|$$

$$\left| B_1^{*} \left( r^* + /\!/ \Phi_t /\!/_{t*} \right) /\!/ P_t^{(1)}\text{-}P_t^{(2)} /\!/_{t^*} 0(a) as\ a \rightarrow 0 \right.$$

*Thus*

$$\left| B_1^{(1)}(a) - B_1^{(2)}(a) \right| \leq (r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ P_t^{(1)} - P_t^{(2)} /\!/_t o(1) as \quad a \rightarrow o$$

And hence

$$\left| \pi_1(P_t^{(1)} - \pi_1\left(P_t^{(2)}\right) \right| \leq ((r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ P_t^{(1)} - P_t^{(2)} /\!/_t o(t) as \quad t \rightarrow o$$

Which together with a similar bound for $\left| \pi_2(P_t^{(1)} - \pi_2\left(P_t^{(2)}\right) \right|$ leads to

$$\| \pi\, (P_t^{(1)} - \pi\, \left(P_t^{(2)}\right) /\!/ \leq (r^* + /\!/ \Phi_t /\!/_{t*}) /\!/ P_t^{(1)} - P_t^{(2)} /\!/_{t*} o(t^*)$$

As $t^* \rightarrow o$ showing that if $t^*$ is small enough, the mapping $\pi$ on

$s_{t^*}(\Phi, r^*)$ is a contraction. Hence the system (2.6)-(2,8) has a unique solution for $o \leq t \leq t^* for\ t^*$ sufficiently small and positive

Now if [0,T] is any finite interval on $R^+$, we can consider [0,T] as a finite union of intervals of length less than or equal to $t^*$ Since the estimate (2.11)is valid for any finite interval we can extend the solution of (2,6)-(2.8) form $[0, t^*]$ to [0,T] and this completes the proof,

3. Stationary age distributions and their linear stability we will now establish the existence of time independent solution

$P_t^*, P_t^*$ Of (1.1)-(1.6): such solutions satisfy

(3.1) $\frac{dp\,P_2^*(a)}{da}+f_i(\mathrm{a}, P_1^*, P_2^*)\,P_i^*(a) = 0,\ \ i = 1,2, a > 0,$

(3.2) $P_i^* = \int_o^\infty P_i^*(a)da, i = 1,\ 2,$

(3,3) $P_i^*=\int_o^\infty b_i(a,\ P_1^*,\ P_2^*)P_i^*(a)da,\ i = 1.\ \ 2$

All solutions of (3.1)are of the from

(3.4) $P_i^*(a) = P_i^*(o)\exp\left[-\int_o^\infty f_i(s, P_1^*, P_2^*)ds\right], i = 1,\quad 2, a \geq 0,$

Which together with (3.3) lead to

$1=\int_o^\infty b_1(a, P_1^*, P_2^*)\exp\left[-\int_o^a f_1(s, P_1^*, P_2^*)ds\right]da,$

(3.5) $1=\int_o^\infty b_2(a, P_1^*, P_2^*)\exp\left[-\int_o^a f_2(s, P_1^*, P_2^*)ds\right]da$

Thus the existence of stationary solutions of (1.1)-(1.6) reduces to the existence of a pair of positive constants $P_1^*, P_2^*$ satisfying (3.5). We can now prove the following,

THEOREM 3. Assume that $f_1$ and $b_i$ satisfy ( $A_1$ ) and ( $A_2$ )and let

(3.6) $F_{i.}\ (P_1, P_2) = \int_o^\infty b_i(a, P_1, P_2)\exp\left[-\int_o^\infty f_i(s, P_1, P_2)ds\right]da,$

$i$ =1, 2, $P_1 \geq 0$, $P_2 \geq 0$,

**And suppose the following hold:**

(3.7) $A_4$. $F_{i.}\ (0,\ 0) > 1,\ i = 1,2$:

$A_5$. *there exists a positive constant c for which*

$F_1$(C,0)< $F_{2.}$(C,0),

(3.8) $F_1$(0.C)< $F_{2.}$(0,C),

Then there exists a unique pair $(P_1^*, P_2^*)$ of real numbers $P_1^* > o$,

$P_2^* > 0$ such that

(3,9) $F_1(P_1^*, P_2^*) = 1 = F_2(P_1^*, P_2^*)$

And the unique nonegative solution of (3.1)-(3.3) is given by

(3.10) $P_2^*(a) = [P_2^* \exp[-f_i(s, P_1^*, P_2^*)ds])/(\int_o^\infty \exp[-\int_o^a f_i(s, P_1^*, P_2^*)ds]da)$,

$i$ =1, 2, $a \geq 0$,

Proof Consider the elements of the surfaces defined by

Z= $F_1(P_1, P_2), z = F_2(P_1, P_2), P_1 \geq 0, P_2 \geq 0$,

In the ($P_1$,$P_2$,z) space. By hypothesis ($A_1$) we have $\partial F_i / \partial P_j < 0$ *and* , by ($A_2$), $\partial b_i / \partial P_j \leq 0$ for all

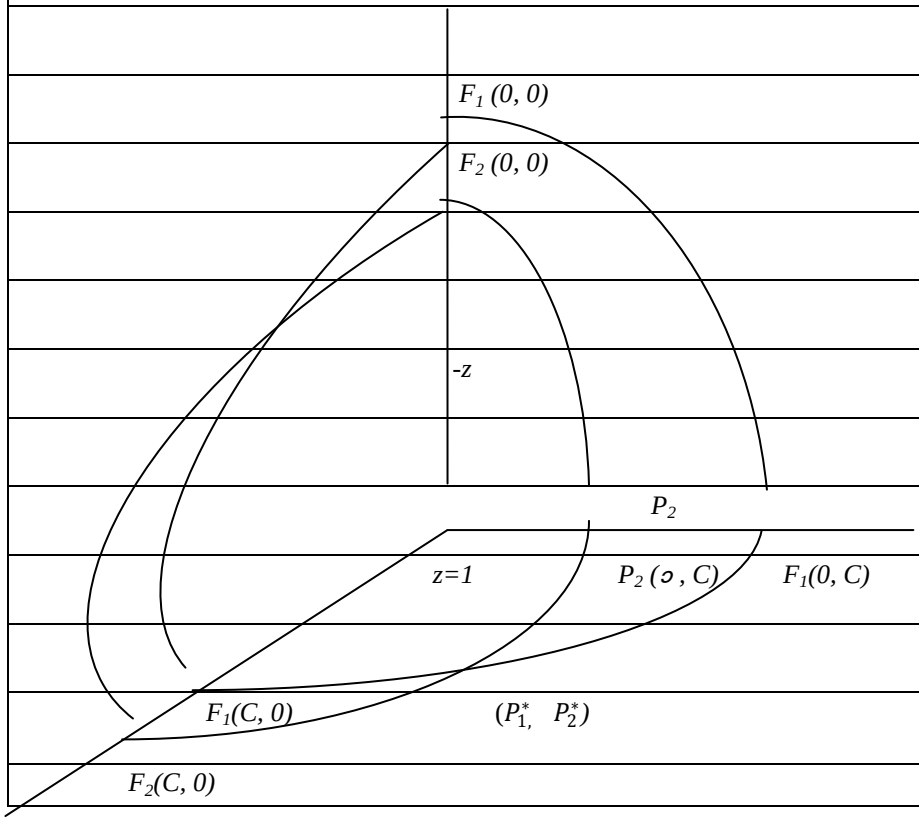
$P_1 \geq 0, P_2 \geq 0$, $i$ =1, 2,

$j$ =1, 2, Hence by ($A_4$) the intersections of these surfaces with the plane z=1 define two curves on the z=1 plane connecting the I ines {z=1, $P_1 \geq 0, P_2 \geq 0,$} and {z= $1, P_2 \geq 0, P_1 \geq 0$, These two curves are defined by

$F_1(P_1, P_2 = 1$ *and* $F_2(P_1, P_2) = 1, F_1(P_1, P_2 \geq 0$,

Which by ($A_s$) will intersect at a unique point say $((P_1^*, P_2^*)$ with $(P_1^* > 0, P_2^* > 0$ (see figure). The uniquess of the point $((P_1^*, P_2^*)$ *is a* consequence of the smoothness of the surfaces $z = F_1(P_1, P_2)$ and $z = F_2(P_1, P_2)$

Following is a hypothetical case for Particles, predications, postulations, ions, quarks, molecules, interactions, transactions, production rates, and destruction rates or mortality rates attributable or ascribable in the transactions and reactionary potentialities.



(The figure corresponds to the case $F_1(0,0) > F_2(0,0) > 0$. The other possibility $1 < F_1(0,0)$ is treated similarly) It will now

Follow that for such a pair $(p^*_1, p^*_2)$ the unique nonnegative solution of (3.1) - (3.3) is given by (3.10).

Now to examine the linear stability of the stationary age distributions $p^*_1, p^*_2$, we let

(3.11)
$$\begin{cases} p_i(a,t) = p^*_i(a) + u_i(a,t) \\ p_i(t) = p^*_i + p_i(t), \quad i = 1, 2, \quad a > 0, \quad t > 0, \\ p_i(t) = \int_0^\infty u_i(a,t)da, \end{cases}$$

In (1.1)-(1.5) and derive the following variational system after neglecting the nonlinear terms in the perturbations $\boldsymbol{u}_i$ and $\boldsymbol{p}_i$ $(i = 1, 2,)$ (such a procedure of linear stability analysis can be justified As has been done in Gurtin and MacCamy [4].)

(3.12)
$$\begin{cases} Du_i(a,t) = -f_i\left(a, p^*_1, p^*_2\right) u_i(a,t) - p^*_i(a)\left\{\sum_{i=1}^{2} \frac{\partial f_i}{\partial p^*_{\square}} p_{i\square}(t)\right\}, \\ u_i(0,t) = \int_0^\infty b_i\left(a, p^*_1, p^*_2\right) u_i(a,t)da + \int_0^\infty p^*_i(a)\left\{\sum_{\square 1}^{2} \frac{\partial b_i}{\partial p^*_{\square}} p_{i\square}(t)\right\}da \end{cases}$$

Where $i$ = 1,2 and (3.13)
$$p_{11}(t) = p_1(t), \quad p_{12}(t) = p_2(t-\tau), \quad p_{21}(t) = p_1(t-\tau), \quad p_{22}(t) = p_2(t)$$

To consider the asymptotic bahaviour as $t \to \infty$ as of solutions of (3.12) we let

(3.14)
$$\begin{cases} u_i(a,t) = u\square(a)\exp[\lambda t], \quad \square(a) \to 0 \text{ as } a \to \infty, \\ p_i(t) = p^*_i \exp[\lambda t], \\ \qquad\qquad i = 1, 2, \quad a > 0, \quad t > 0, \\ p^*_i = \int_0^\infty \square(a)da, \end{cases}$$

Using (3.14) in (3.12) - (3.13),

$$\frac{d\square}{da} + \lambda\square = -f_i\left(a, p^*_1, p^*_2\right)\square(a) - p^*_i(a)\left\{\sum_{j=1}^{2} \frac{\partial f_i}{\partial p^*_j} p^*_{i\square}\right\},$$

(3.15)

$$\mu_i(0) = \int_0^\infty b_i\left(a, p_1^*, p_2^*\right)\mu_i(a)da + \int_0^\infty p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial b_i}{\partial p_j^*}\ p_i^*\right\}da,$$

Where $i = 1,\ 2,$ and

$$p_{11}^* = p_1^*, \quad p_{12}^* = p_2^* \exp\{-\lambda\tau\},$$

$$p_{21}^* = p_1^*,\ \exp\{-\lambda\tau\},\ p_{22}^* = p_2^*.$$

If we integrate the first of (3.15) and use the fact that $\mu_i(a) \to 0$ as $a \to \infty$ we can derive that $p_1^*$ and $p_2^*$ are governed by

(3.16) $$\lambda p_1^* - \int_0^\infty p_1^*(a)\left[\left(\frac{\partial b_1}{\partial p_j^*} - \frac{\partial f_1}{\partial p_1^*}\right)p_1^* + \left(\frac{\partial b_1}{\partial p_2^*} - \frac{\partial f_1}{\partial p_2^*}\right)p_2^* e^{-\lambda\tau}\right]da$$

$$= \int_0^\infty \left[b_1\left(a, p_1^*, p_2^*\right) - f_1\left(a, p_1^*, p_2^*\right)\right]\mu_1(a)da,$$

(3.17) $$\lambda p_2^* - \int_0^\infty p_2(a)\left[\left(\frac{\partial b_2}{\partial p_1^*} - \frac{\partial f_2}{\partial p_1^*}\right)p_1^* e^{-\lambda\tau} + \left(\frac{\partial b_2}{\partial p_2^*} - \frac{\partial f_2}{\partial p_2^*}\right)p_2^*\right]da$$

$$= \int_0^\infty \left[b_2\left(a, p_1^*, p_2^*\right) - f_2\left(a, p_1^*, p_2^*\right)\right]\mu_2(a)da,$$

From the first of (3.15) we get

(3.18) $$\mu_i(a) = \left\{\mu_i(0)p_i^*(a)/p_i^*(0)\exp[\lambda a]\right\} - p_i^*(a)\exp[-\lambda a]\int_0^a \left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right\}e^{\lambda s}ds, \quad i = 1, 2,$$

Which by the second of (3.15) becomes,

(3.19) $$\mu_i(a)\left[p_i^*(a)/p_i^*(0)c_i\exp(\lambda a)\right]\left[\int_0^a p_i^*(a)\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right\}da\right.$$

$$\left. - \int_0^\infty b_i\left(a, p_1^*, p_2^*\right)p_i^*(a)\,e^{-\lambda a}\left\{\int_0^a\left(\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right)e^{\lambda s}ds\right\}da\right]$$

$$- p_i^*(a)\,e^{-\lambda a}\int_0^a\left\{\sum_{j=1}^{2}\frac{\partial f_i}{\partial p_j^*}\ p_i^*\right\}e^{\lambda s}ds, i = 1, 2,$$

where

(3.20) $$c_i = 11 - \frac{1}{p_i^*(0)} \int_0^\infty b_i\left(a, p_1^*, p_2^*\right) p_i^*(a) e^{-\lambda a} da, \quad i = 1, 2,$$

Now supplying $\varphi_1(a)$ and $\varphi_2(a)$ from (3.19) in (3.16)-(3.17) we derive that

$$\left(\lambda - A_{11}\right) p_1^* - A_{12} p_2^* \exp\left(-\lambda\tau\right) = g_{11}(\lambda) p_1^* + g_{12}(\lambda) p_2^* \exp(-\lambda\tau),$$

(3.21)

$$-A_{21} p_1^* \exp\left(-\lambda\tau\right) + \left(\lambda - A_{22}\right) p_2^* = g_{21}(\lambda) p_1^* \exp\left(-\lambda\tau\right) + g_{22}(\lambda) p_2^*$$

Where

(3.22) $$A_{i\varphi} = \int_0^\infty p_i^*(a) \left[\frac{\partial b_i}{\partial p_\varphi^*} - \frac{\partial f_i}{\partial p_\varphi^*}\right] da,$$

(3.23) $$g_{i\varphi}(\lambda) = \left[p_i^*(0) c_i\right]^{-1} \left\{\int_0^\infty \sigma_i(a) p_i^*(a) e^{-\lambda a} da\right\} \varphi \left\{\int_0^\infty \left[p_i^*(a) \frac{\partial b_i}{\partial p_\varphi^*} - b_i\left(a, p_1^*, p_2^*\right) p_i^*(a) e^{-\lambda a} \left(\int_0^a \frac{\partial f_i}{\partial p_\varphi^*} e^{\lambda s} ds\right)\right] da\right\}$$

$$- \int_0^\infty \sigma_i(a) p_i^*(a) e^{-\lambda a} \left(\int_0^a \frac{\partial f_i}{\partial p_\varphi^*} e^{\lambda s} ds\right) da, \quad i, \varphi = 1, 2,$$

(3.24) $$\sigma_i(a) = b_i\left(a, p_1^*, p_2^*\right) - f_i\left(a, p_1^*, p_2^*\right), \quad i, = 1, 2.$$

To solve the linear perturbational system (3.12) it is enough to find the constants $\lambda, p_1^*, p_2^*$ satisfying (3.21) since one can then use (3.19) to find the perturbations from (3.14). A nontrivial solution $\left(p_1^*, p_2^*\right)$ for (3.21) will exist if and only if $\lambda$ is a root of the equation $et \begin{bmatrix} \lambda - \left(A_{11} + g_{11}(\lambda)\right) & -\left(A_{12} + g_{12}\right) e^{-\lambda\tau} \\ -\left(A_{21} + g_{21}\right) e^{-\lambda\tau} & \lambda - \left(A_{22} + g_{22}(\lambda)\right) \end{bmatrix} = 0$

Or equivalently

(3.25) $$\lambda^2 - \lambda\left(A_{11} + A_{22}\right) + A_{11}A_{22} - A_{12}A_{21} \exp\left[-2\lambda\tau\right] = \lambda s_1(\lambda) + s_2(\lambda, \tau)$$

Where $S_1(\lambda) = g_{11}(\lambda) + g_{22}(\lambda)$ ,

$S_2(\lambda, \tau) = \left[A_{21} g_{12}(\lambda) + A_{12} g_{21}(\lambda) + g_{12}(\lambda) g_{21}(\lambda)\right] \exp\left[-2\lambda\tau\right]$

$$-A_{11} g_{22}(\lambda) - A_{22} g_{11}(\lambda) - g_{11}(\lambda) g_{22}(\lambda)$$

The following result establishes the linear local asymptotic Stability of the stationary age distributions $p_1^*(a)$ and $p_2^*(a)$

THEOREM 4 Assume the following:

$A_6.\ 0 < \beta_i^* < d_i^*\ (i,\ =1,\ 2.)\ \ (see(A_1)\ and\ (A_2));$

$A_7.\ \{A_{11} + A_{22} + \| S_1(0) \|\}^2 < 4\{(A_{11}A_{22} - A_{12}A_{21}) - \| S_2(0) \|\},$

$$\| S_1(0) \| = |g_{11}(0)| + |g_{22}(0)|,$$

$$\| S_2(0) \| = |A_{11}\, g_{12}(0)| + |A_{12}\, g_{21}(0)| + |g_{12}(0)\, g_{21}(0)|$$

$$+ |A_{11}\, g_{22}(0)| + |A_{22}\, g_{11}(0)| + |g_{11}(0)\, g_{22}(0)|$$

Then all the roots of (3.25) have negative real parts and hence the

perturbations $u_i\ (a,\ t) = \phi(a) \exp[\lambda t] \to 0\ as\ \ t \to \infty$ for $i = 1, 2\ and\ a \geq 0$

Proof . Define $F_1(\lambda, \tau)$ and $F_2(\lambda, \tau)$ as follows:

$F_1(\lambda, \tau) = \lambda^2 - \lambda(A_{11} + A_{22}) + A_{11}A_{22} - A_{12}A_{21} \exp[-2\lambda\tau];$

$F_2(\lambda, \tau) = \lambda S_1(\lambda) + S_2(\lambda, \tau),$

Using $(A_6)$ one can show that for $\lambda$ with re $(\lambda) \geq 0$, $c_1$ and $c_2$ are positive and bounded away from zero, Now since $A_{ij} > 0$ and $\tau \geq 0$ we find from the nature of the dependence of $g_{ij}$ on $\lambda$ that for all $\lambda\ \ \mu \pm iw$ with $\mu \geq 0$ and $w \geq 0$ we have

(3.26) $$\left| F_1(\lambda, \tau) - P_2(\lambda, \tau) \right|$$

$$\geq \left| F_1(\lambda, 0) | - F_2(\lambda, \tau) \right|$$

$$\geq |\lambda|^2 - |\lambda|\{A_{11} + A_{22} + \| S_1(0) \|\} + \{A_{11}A_{22} - A_{12}A_{21} - \| S_2(0, \tau) \|\}$$

$$\geq \left[(|\lambda| - P)^2 + (q)\right]^2$$

Where

$p = \{A_{11} + A_{22} + \| S_1(0) \|\}/2,$

$q = \left\{4 | (A_{11}A_{22} - A_{12}A_{21}) - \| S_2(0, T) \| | - | A_{11} + A_{22} + \| S_1(0) \| |^2\right\}/2,$

It will now follow, from (A) and (3.26), $F_1\ (\lambda, \tau) = F_2(\lambda, \tau)$ connot have roots with zero or positive real parts and this completes the proof.

In conclusion we remark that the condition (3.7) means that the net creation rates of each system is greater than utility; a condition of this type is well known in age wise classification of particles, atoms, molecules, cosmological phenomenon. (3.8) can be interpreted to mean that the intraspecific inhibition is higher than that of the

interspecific interaction. The conditions of Theorem 4 are analytical and do not lend themselves for any worthwhile interpretation in terms of the model parameters. So also an interpretation can be recorded as a remark in respect of destruction rates of particles, properties, elements, or projections.

Note: The model is applicable to each and every module. As stated in the foregoing, there are possibilities of situation where an account itself be a receiver and it is also the one to which debit is raises. Such examples are legion in all the fields of finance, economics, political science and physics and chemistry, and of course mathematics. Classification based on time is the preeminently and predominantly suited premise for such model. We repeat that the model is to be read with each module We will expatiate the importance and implicational ramifications as we proceed with the work. Assortment-blend, collection-combinations, combo-gallimaufry, goulash-hash, hodgepodge-medley, mishmash-mixture, motley-pastiche patchwork-salmagundi, soup-stew, mixed bag-mélange, olio- olio podrida, assortment variety, Array choice, collection combination, combo diversity, hodgepodge jumble, kind medley, miscellany mishmash, mixed bag mixture, mélange potpourri, selection sort, brew concoction, beverage blend, broth compound, distillation drink, fermentation hash, hodgepodge infusion, instillation liquor, mélange miscellany, mishmash mixture, mélange potpourri, borscht bouillon, bowl brew, chowder concoction, consommé decoction, gumbo hodge-podge, olio porridge, potage potpourri, pottage puree, splash stock, vichyssoise hodgepodge, mixture- mess, collection combination, goulash hash, jumble medley, miscellany mishmash, mixed bag mélange, olio patchwork, potpourri salmagundi, jumble clutter, disarrangement disarray, disorder farrago, gallimaufry garbage, goulash hash, litter medley, mess miscellany, mishmash mixture, muddle mélange, olio pastiche, patchwork potpourri, salmagundi scramble, shuffle snarl, tangle tumble, medley miscellany, assortment brew, collection combo, conglomeration farrago, hodgepodge jumble, melee mingling, mishmash mixed bag, mixture mélange, pasticcio pastiche, patchwork potpourri, salmagundi variety, mélange mixture, assortment combo, farrago gallimaufry, pasticcio pastiche, patchwork potpourri, salmagundi soup, stew miscellany, accumulation aggregation, anthology assortment, brew collectanea, combination combo, compilation conglomeration, accumulation diversity, farrago gallimaufry, melee mishmash, mix mixed bag, mixture muddle, mélange odds and ends, olio pasticcio, pastiche patchwork, potpourri salad, salmagundi smorgasbord, stew variety, mishmash jumble, medley mess, miscellany mixed bag, mixture mélange, olio patchwork, potpourri salmagundi, mixed bag array, assortment choice, collection combination, combo conglomerate, diversity hodgepodge, jumble medley, mélange miscellany, mishmash mixture, montage mélange, odds and ends, patchwork potpourri, smorgasbord sundries, mix/mixture assortment, combination, admixture adulteration, alloy amalgam, amalgamation assimilation, association batter, blend brew, combine combo, commixture composite, compound concoction, confection conglomeration, cross crossing dough fusion, goulash grab bag, hodgepodge hybrid, hybridization incorporation, infiltration interfusion, jumble mash, medley merger, mingling miscellany, mishmash mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, odds and ends assortment, bits and pieces, debris etcetera, hodgepodge jumble, knickknacks leavings, leftovers litter, medley melee, miscellaneous paraphernalia, miscellany mishmash, mixed bag motley, mélange notions, novelties oddments, olio particles, patchwork potpourri, rags refuse, remainder remains, remnants rest, rummage scraps, sundries sundry items, compilation mishmash, paste-up patchwork, potpourri reappropriation, reproduction synthesis, conglomeration cross, crossing dough, fusion goulash, interfusion jumble, mash medley, merger mingling, miscellany mishmash, mix mixed bag, mosaic mélange, package patchwork, potpourri salmagundi, saturation soup, stew transfusion, union variety, colluviums gallimaufry, hash mixture, pasticcio pastiche, conglomeration cross section, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multiformity multiplicity, mélange potpourri, range shift, variegation variousness, divergencies diverseness, departure discrepancy, disparateness divergencies, diverseness diversification, diversity fluctuation, heterogeneity heterogeneousness, incongruity intermixture, many-sidedness medley, miscellaneousness miscellany, mishmash mixed bag, mixture modification, multifariousness multiformity,

mélange potpourri, range shift, soup stew, variance variation, variegation variousness, oddment odds and ends, assortment bits, polymorphism assortment, change collection, combo conglomeration, cross section departure, discrepancy disparateness, divergencies diverseness, diversification diversity, fluctuation heterogeneity, heterogeneousness incongruity, intermixture many-sidedness, medley miscellaneousness, miscellany mishmash, mixed bag mixture, notions novelties, oddment oddments, olio particles, patchwork potpourri of the particles, interactions, accounts, transactions in the cosmos. The wisdom and understanding of our own divine essence is like a beaconing begorra, bureaucratic burdock, befuddling binary and digital digamma, that invites us inward toward ever deeper connection with all around us. Yet rather than look within, most in this world prefer to look outward for divine wisdom, thus sustaining a dependence upon a vast hierarchy that stretches between the individual and the Divine. In all our wanderings away from the Divine, humankind has obscured, obfuscated, obliterated, adjured, abnegated and abjured its most compelling features through a persistent belief systems in limitations arising out of the controlling mechanisms of this hierarchy of sub ordinates, superordinates and coordinates. Has my consciousness changed? Why that doubt? Twist and shake, ponderous novelty themes, pathos, poignancy in writing in different styles, accord-coetaneousness, coevality-coincidence, concurrence-conformity, conjunction-contemporaneousness, harmony, order, peace, simultaneousness- synchronicity exciting optimism of youth, power obmbudsman,pulsations and throbbing of pop, hard luck stories of people, modernity and mischief, grimalkins of rurality, wobbling boats of lovers, saxophone in mouths all point to newer value orientations. The Divine dances and the “Grand design “is perpetrated outside of the confines of any hierarchical structures, be it feudal, organizational. It is complete within itself, and has a singular purpose of demonstrating the collective potential of all life within the universe. It is the archetype of perfection. It is the standard bearer of each soul's innate design and ultimate destiny. The essence of the Divine is far beyond mental conception, yet humanity's tendency is to resort to the limiting language of the hierarchical paradigm acceding- accession, accommodation- accord, accordance, acknowledging, adjustment-affiliation, affinity-alliance, amity-approving, arbitration-arrangement, assenting, authorizing, bargaining-compatibility, compliance--compromise, concert-concession, conformity-congruity, consistency-correspondence, endorsing, granting-harmony in conformity to pen chance ,predilections of the protagonist to define and understand it. Searching for Connection and Wholeness calls for not lumping elements like me made of clay feet and weak knees, owing allegiance to the fringe, without courante and courage of conviction and will for vindication, visual restlessness and exaggerated narratives of self to its own, exciting potentialities available, terrestrial passion, and skew of scams all bear an infallible demonstration for new order of consciousness. Theorists have posited a virtual smorgasbord of objects to account for dark matter, although many of them have fallen prey to observational constraints. As leading possible candidates for baryonic dark matter, there are black holes (large and small), brown dwarfs (stars too cold and faint to radiate), sun-size MACHOs, cold gas, dark galaxies and dark clusters, to name only a few. The range of particles that could constitute nonbaryonic dark matter is limited only slightly by theorists' imaginations. The particles include photinos, neutrinos, gravitinos, axions and magnetic monopoles, among many others. Of these, researchers have detected only neutrinos--and whether neutrinos have any mass remains unknown. Experiments are under way to detect other exotic particles. If they exist, and if one has a mass in the correct range, then that particle might pervade the universe and constitute dark matter. But these are very large "ifs." To a great extent, the details of the evolution of galaxies and clusters depend on properties of dark matter. Without knowing those properties, it is difficult to explain how galaxies evolved into the structures observed today. As knowledge of the early universe deepens, I remain optimistic that we will soon know much more about both galaxy formation and dark matter. What we fail to see with our eyes, or detectors, we can occasionally see with our minds, aided by computer graphics. Computers now play a key role in the search for dark matter. Historically, astronomers have focused on observations; now the field has evolved into an experimental science. Today's astronomical experimenters sit neither at lab benches nor at telescopes but at computer terminals. They scrutinize cosmic simulations in which tens of thousands of points, representing stars, gas and dark matter; interact gravitationally over a galaxy's lifetime. A cosmologist can tweak a simulation by adjusting the parameters of dark matter and then watch what happens as virtual galaxies evolve in isolation or in a more realistic, crowded universe.

Old consciousness will not do. Newer one has to be imbibed and cultivated. When people are unaware of their inner wisdom and wholeness, they tend to search for order and security outside themselves. That is my personal experience. There is always search for sanctity of emotion, aesthetic expression, pursuit of harmony, unraveling realms, syncopated ending and stillness of mind and shiver of heart at the mirror broken somewhere are the strains of newer wisdom if it can be called one. Same questions shall be asked and same shall be answer in the similar vein. Uncertain of their place within the hierarchical order of the world, they define themselves based upon their insecurities. Individuals thus become only pieces of their wholeness and like shards of glass from a beautiful vase; they bear little resemblance to their aggregate beauty. Within the hierarchy, many in high positions of power have taken advantage of our collective insecurities in attempting to guide the development of all humankind. They have obscured the direct connection between the individual and the Divine through a variety of means designed to intercede between our inner essence and our divine source. Each individual can come to know them to be free of all forms of institutional and organisational and societal hierarchical control. This is not to imply that we should not trust others or join together in bonds of friendship and community. It is simply a reminder that relative truth is constantly shifting in the hands of those who desire to control. And even when the motive for their controlling behaviors may be of good will, it is still a form of control. When the revealers of "truth" within the hierarchy withhold and suppress information, they are usually positioning themselves to maintain or acquire more power rather than to disseminate empowerment to all. The desire for connection and wholeness is a fuel that drives us to seek out and explore the hierarchical paradigm. This inner longing provides us with the motivation to seek help and guidance from a specific group within the hierarchy, and in so doing, cultivate a sense of belonging and connection. Furthermore, the hierarchical paradigm is a stage whereby we develop a sense of connection to some grand, encompassing vision. This is why the hierarchy nurtures prophets who point toward a greater vision. Spiritual leaders are able to peer deeply beneath the surface reality of life and experience how intricately connected every life form is, and how the composite of all life is intelligent far beyond measure. These visionary leaders can thus interpret reality through their personal abilities to perceive and express life's dimensional depth and limitless intelligence. Yet no one is able to articulate life's full dimensional depth and breadth with the tools of language. They can only, at best, describe their interpretation or their impressions. In actuality, all of us are able to peer beneath the surface reality of life and perceive a unique vision of the universe. We require only time and intention to develop our own interpretations. And this is precisely what many great spiritual leaders have taught. Life's deeper essence is not an absolute to be experienced by the chosen few, but an evolving, dynamic intelligence that wears as many faces as there are life forms. No group or people have the exclusive portal into the universe by which the Divine expresses itself in all its majesty. The portal is open and available to all, because the Divine is within all things. Those recognized as great prophets each produced a vision of the universe beyond what was currently defined by the hierarchy. Because their interpretations were articulated with authority and depth of insight, they became a target of debate among various groups in the hierarchy. This debate then created a polarity of belief. A sympathetic constituency emerged to defend and embellish their leader's interpretation, while established groups held it in contempt of previously held beliefs. Invariably, the leader's vision became confined and shaped into dogma by followers who desired to create a new religion or sect. Thus, this infusion of fresh insight quietly receded into the hands of the hierarchy, where its deeper meaning was obscured by the very fact that it was incorporated into a rigid structure that both protected and promoted it. A new paradigm is emerging that promotes a clear connection of individual consciousness to the compelling features of the Divine without the intervention of a hierarchy. This is when the fables and myths of history step into the light and become known as they were originally intended. This is the time when language will be transformed into a new form of communication that breaks down all barriers of control. Personal transformation, through the awakening of inner wisdom and understanding, is the pathway into wholeness. The transformation paradigm is initiated simply by the recognition that rather than the dependency-inducing ways of the hierarchical paradigm, there are accelerated, independent pathways that bypass the hierarchy and lead to self-mastery. These new pathways lead to the divine wisdom and understanding that is present within all of us. This wisdom can be accessed through the practice of three principles of transformation: seeing the Divine in

all, nurturance of life, and gratitude. The application of these life principles disengages individuals from the controlling elements of the hierarchy, thus initiating the transformation experience. As there are relative truths, there are relative freedoms. As individuals evolve through the hierarchical process, an ever-increasing sense of freedom is gained, yet external forces continue to exert control through limiting language, confining belief systems, and more. These controlling influences lead to anatomical architecture, body, build, cagey carcass, construction, enclosure, fabric, flounce, form, framework, fringe, groundwork, hem, mount, mounting outline, physique, scaffold, scaffolding scheme, setting shell, stage, structure, support, system, trim, trimming, truss- valance continued reliance upon the hierarchy as it unceasingly attempts to impose a sense of inequality between us. The underlying equation of the hierarchical paradigm is: individual + hierarchy = God connection. In the case of personal transformation it is: individual + inner wisdom and understanding divine equality with all. The time has come to integrate the dominant paradigm of the hierarchy with the liberating transformation paradigm. That shall redeem all promises made to oneself. There shall be a Matrix force or balancing force that shall remain UN satisfies until conservativeness is reached. This integration occurs naturally once we have fully explored the two paradigms and develop a synthesis paradigm whereby transformation is attained by searching for connection and wholeness through our own inner wisdom and understanding. It is this combination of release from dependency on the hierarchy and transformation from within that initiates the synthesis paradigm. Once we take responsibility for our own transformation and integration, it does not mean that the hierarchy is to be shunned or avoided. The hierarchy is quite benign as a manipulative force and merely represents one important stage in the journey toward wholeness. What is being set in motion now is the initial preparation for these paradigm shifts. More specifically, these paradigms will be simultaneously played out over the coming years. As always, it will be the choice of each one of us as to which paradigm we embrace in our journey. All of the highest imaginings of humanity are yet unaware of our deepest foundation. We have sought the upper reaches of the building, yet remain unaware of the foundation's design. It is here, at the very core of existence that the Divine is bursting forth with its creative energy while simultaneously reintegrating with its invitation to wholeness. It is here that equality is realized, not in the lofty places of relative truth lodged in the hierarchy, but rather in the deepest part of the foundational plan of life's original source and ultimate destiny. The origin and destiny of existence is the tone of equality in life. Listen for this tone – this frequency of vibration – and follow it back into the very foundation from whence all things arise and return. Consider these words as symbols only. Feel the truth that stands behind these words, and tap into this empowering energy force that reaches out for you. Know it as a tone or vibration – a resonance that waits for you around every corner in which your life will turn. It is a beacon of the Divine gathering itself into the form of language in order to usher you to a place from which you can experience the formless tone of equality – the bypass of limitation. It is the primal language of our divine source that bestows to you the freedom to generate your own deepest beauty in the expression of your highest truth.(Reference: Wikipedia)

## Gopalasamy Model For Procrastinated Conducive, Felicitous Counter Sallies In Two Specified Systems

Abstract

Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of decay rate of perturbations. Sole intention, main objective and primary concern of

the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are counted in the same vein as that of predator prey. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 15. Preliminary remarks

Almost all the equations of population ecology have been derived starting from the following simple format (Note this can be used for the development of the creed and type of systemic types in case of cosmology and astrophysics and fiancé and econometrics. To give an example of finance, we can think of a account of the company giving rise to various accounts of its employees, and that is the contribution of the single account for the development of the total system types of the given configuration). It may be account, particle, atom, molecule, proposition, problem or any such thing. We shall delineate and disseminate information and particulars as and when the model shall be applied to the problem in question.

$$\frac{dN(t)}{dt} = \text{an individual's contribution}$$

*to population growth* $N(t)$

(1.1)

where $N(t)$ denotes the size of a population of a single species (or biomass) with overlapping generations at time $t$, and $dN(t)/dt$ is the rate of change of the population size at time $t$. Subsequently one makes an assumption regarding the factor in the brackets in (1.1); in particular if one assumes that an individual entity's contribution to the population growth is denoted by a function say, f(t, N), defined suitably for all t>0, N then one obtains from (1.1) the so called Kolmogorov formulation in the form

$$\frac{dN(t)}{dt} = f(t, N(t))\, N(t).$$

(1.2)

Various choices of $f$ together with some ecologically plausible assumptions such as the temporal constancy of the environment and density dependent effects in $f$ lead to various equations. For instance if $f(t, N) \equiv r$ (a positive constant) one obtains the Malthusian formulation in the form

$$\frac{dN(t)}{dt} = rN(t)$$

And if one assumes $f(t, N) \equiv r - (r/K) N$ for some positive constants r and K one gets the familiar logistic equation of population biology

$$\frac{dN(t)}{dt} = rN(t) \left\{ 1 - \frac{N(t)}{K} \right\}$$

In order to device models with time delayed effects, such effects are introduced in the average growth rate $F(t, N)$ of (1.2); for instance if $f(t, N) \equiv r - (r/K) N(t - \tau)$ for some constant $\tau > 0$ we obtain Hutchinson's[19]delay loistic equation

$$\left\{ \frac{dN(t)}{dt} = rN(t) \left\{ 1 - \frac{N(t-\tau)}{K} \right\} \right. .$$

(1.3)

If for instance

$$F(t, N) = r - (r/K) N(t) - \int_{-\infty}^{t} k(t - s) F(N(s))\, ds$$

Where $k$ is a piecewise continuous nonnegative valued function such that $\int_0^\infty k(s)\,ds = 1$ with some appropriate function $F$, we will get Volterra's model of a population which pollutes its environment and pollution has an accumulative and toxic effect; such an equation is of the form

$$\frac{dN(t)}{dt} = N(t) = \left\{ r - (r/K) N(t) - \int_{-\infty}^{t} k(t-s) F(N(s))\,ds \right\} \quad (1.4)$$

Which has been studied in different formats by several authors (Cushing [12], Cohen *et al.* [10], and Gopalasamy [16])

One of the significant advantages of deriving various models from the prototype (1.2) is that if the initial; conditions (initial population sizes) are nonnegative then the nonnegativity of the population size $N(t)$ for $t > 0$ follows from the fact that solutions of (1.2) satisfy

$$N(t) = N(0) \exp\left[\int_0^t f(s, N(s))\,ds\right]$$

Which implies that $N(t) \geq 0$ if $N(0) \geq 0$.

While the time delayed logistic equation (1.3) explains some of the observed oscillations of certain populations, its derivation has not been entirely satisfactory. For instance it has been pointed out by Ricklefs ([27], page 488) that "time lags might be expected to occur primarily in stabilizing processes involving production rather than death because death is an immediate response to environmental change". Recently Cushing [13] has examined the effects of response and maturation delays in a number of single and multi system models; as it has been noted by Cushing [13], the models involving delayed reproduction or not necessarily derivable from the Kolmogorov type systems of the form (1.2) by suitably defining the average growth rate $(1/N)(dN/dt) = f(t, N(t))$. While Cushing [13] has been mostly concerned with the possibility of delay induced oscillations we will consider a class of delays in reproduction and interspecific interaction which maintain the stability of the systems.

Our starting point to drive models with delayed reproduction is the following balance equation which assumes that there is no immigration or emigration. Account, particle, atom, molecule, proposition, problem et al

$$\frac{dN(t)}{dt} = \text{birth rate} - \text{death rate}. \quad (1.5)$$

Now for instance if we consider a population of adult flies then the birth rate at time $t$ depends on the adult population at time $t^{-}\ \tau$ where $\tau$ is the time required for the larvae to become adult flies. If the birth and death rates are governed by density dependent factors then we have from (1.5) the following time delayed model

$$\frac{dN(t)}{dt} = b\big(N(t^{-}\tau)\big) - m(N(t))$$

(1.6)

Where the functions $b$ and $m$ denote the density dependent birth and death rates of Account, particle, atom, molecule, proposition, problem et al respectively. If the time delay is continuously distributed then one can consider in the place of (1.6) an equation of the form

$$\frac{dN(t)}{dt} = b\int_{-\infty}^{t} k\{(t-s)N(s)ds - m(N(t))$$

Or equivalently

$$\left\{\frac{dN(t)}{dt} = b\int_{0}^{\infty} k(s)\right\}N(t-s)ds - m(N(t))$$

(1.7)

Where the delay kernel $k$ denotes the weight of the delayed of the delayed effects.

It is not always obvious that suitable nonnegative initial conditions for (1.6) and (1.7) will imply the nonnegativity of the solutions of (1.6) and (1.7). The purpose of this article is to show that in two species systems of competitive and prey- predator communities, delayed reproduction and delayed interactions donot destabilise the systems if the self- regulating negative feedback effects are sufficiently strong compared to interspecific interactive effects and if the self regulating effects are realised with no time delays. When we say that time delays do not destabilise, we will mean that delays do not render an asymptotically (locally) stable steady state unstable in such a manner that the loss of stability either leads

To a delay – induced bifurcation to persistent and undamped oscillations (as in the case of Hopf- bifurcation) or the relevant linear variational system has unbounded solutions; that is a mathematical instability is not induced by delays in the sense, that the roots of the characteristic equation are pushed to the imaginary axis or to the right of it on the complex plane. However delays can make an otherwise stable system less stable which means that the rate of convergence to zero of solutions of the variational system can decrease with an increase of delay and the variational system can even oscillate (perhaps violently) before converging to the trivial state. Such a reduction in stability can

happen even when all the roots of the characteristic equation have negative real parts. For our purposes such a reduced stability is still stability. In order to see the effects of delays on such less stable systems we will propose below a method of estimating the rate of decay of local perturbations based on the variational system and indirectly relying on the characteristic equation. For a special class of two species models with a single delay we have obtained (Gopalsamy [16]) estimates of decay rates (or return times) usingthe characteristic equation; such a procedure is almost impossible in the case of many species systems with several delays. The method we propose in Section 6 below for the estimate of decay rates can be adapted to *n*- species systems with $n^2$ discrete delays.

Let $x_1(t)$ and $x_2(t)$ denote the population sizes (or biomasses) of two systems competing for a common pool of resources with common pool of resources.. Let $b_i$ and $m_i$ $(i = 1, 2)$ denote the respective density dependent birth (potential) and death rates so that in the absence of any time delays the population sizes are governed by (of Account, particle, atom, molecule, proposition, problem et al). In addition as usual, we assume that the system studied follow Haimovici's equations, for which the dove tailing explanation has been appended

$$\frac{dx_1(t)}{dt} = b_1\left(x_1(t)\right) - m_1(x_1(t)), x_2(t))$$

$$\frac{dx_2(t)}{dt} = b_2\left(x_2(t)\right) - m_2(x_1(t)), x_2(t))$$

(2.1)

In order to make the system (2.1) represent a competition system we will make the following assumptions on the birth and death rates; $b_i$, $m_i$ $(i = 1, 2)$ are continuous with continuous partial derivatives for all $x_i \geq 0$ $(I = 1, 2)$; also we assume (see for instance Rescigno and Richardson [26]) $\frac{\partial b_i}{\partial x_1} > 0, \frac{\partial m_1}{\partial x_j} > 0$ for $x_i > 0$ $(i, j = 1, 2)$, (2.2)

(xxii) $b_1(0) = 0$, $m_1(0, x_2) \equiv 0, b_2(0) = 0$, $m_2(x_1, 0) \equiv 0$,

(2.3)For some $x_{1,}^* > 0, x_{2,}^* > 0$ we have

$$b_1(x_1^*, -m_1(x_1^*, 0) = 0,$$

$$b_2(x_2^*, -m_1(0, x_2^*) = 0,$$

(2.4)

(xxiii) There exist positive constants $\delta_1$, $\delta_2$ such that

$b_1(\delta_1) - m_1(\delta_1, x_2) < 0$,

$b_2(\delta_2) - m_2(x_1, \delta_2) < 0$, for $x_1 \geq 0$, $x_2 \geq 0$,

(2.5)

(xxiv) There exists a point $(\alpha, \beta), \alpha > 0, \beta > 0$ such that

$$b_1(\alpha) - m_1(\alpha, \beta) = 0 = b_2(\beta) - m_2(\alpha, \beta).$$

(2.6)

The conditions on $b_i$ in (2.2) mean that the potential (not average) birth rates are density dependent and any negative density dependent feedbacks to the opened or birth rate of Account, particle, atom, molecule, proposition, problem et al rates are included in the destruction rates; the conditions on $m_i$ indicate intraspecific and interspecific confederation needs and necessities. (2.3) implies that (0, 0) is a (trivial) steady state of the system (2.1) as it is in the case of all ecological models; (2.4) means that in the absence of either of the two species the other one has a nontrivial steady state; (2.5) implies that each of the two populations cannot grow unboundedly since when $x_1(t) > \delta_i, \frac{dx_1}{dt} < 0$. (2.6) means that there exists a nontrivial steady state $(\alpha, \beta)$ in the interior of the positive quadrant of the $(x_1, x_2)$- plane. These assumptions are standard and are minimal in the opinion of the author.

An immediate example of (2.1) –(2.6) is the familiar Volterra –Lotka competition system described by

$$\frac{dx_1(t)}{dt} = r_1 x_1(t) - a_{11} x_1^2(t) - a_{12} x_1(t) x_2(t),$$

$$\frac{dx_2(t)}{dt} = r_2 x_2(t) - a_{21} x_1(t) x_2(t) - a_{22} x_2^2(t),$$

(2.7)

Where $r, a_{ij}, (i, j = 1,2)$ are positive constants satisfying the condition

$$\frac{a_{11}}{a_{21}} > \frac{r_1}{r_2} > \frac{a_{12}}{a_{22}}.$$

One of the first questions to be resolved for (2.1) is the following: if $x_i(0) > 0$, $x_2(0) > 0$ does it follow that the corresponding solutions of (2.1) satisfy the condition $x_1(t) \geq 0$, $x_2(t) \geq 0$ for those $t$ for which a solution of (2.1) exists? In the $(x_1, x_2)$ phase plane of (2.1) we have $(0, 0), (x_1^*, 0), (0, x_2^*), (\alpha, \beta)$ as the steady states of (2.1) and also it will follow from (2.3) that the co ordinate axes

$x_1 = 0$ and $x_2 = 0$ are invariant sets f (2.1) ; such invariance of the coordinate axes will imply that if $x_1(0) > 0$, $x_2(0) > 0$ then the corresponding population trajectory $(x_1(t), x_2(t))$ of (2.1) cannot reach the outside of the nonnegative quadrant $x_1 \geq 0$, $x_2 \geq 0$.

Our plan is to derive first a set of sufficient conditions for the coexistence of the two competing species and then in the next section show that under the same conditions the two species can coexist if there are time delays in the reproduction and interspecific interactions. We recall that the two competing species described by (2.1) are said to coexist if and only if the nontrivial steady state $(\alpha, \beta)$ is (at least locally) asymptotically stable. To examine such a stability of $(\alpha, \beta)$ we derive the corresponding linear variational system for the perturbations $X_1, X_2$ where

$$x_1(t) \equiv \alpha + X_1(t), \qquad x_2(t) \equiv \beta + X_2(t)$$

Such a variational system is found to be

$$\left.\begin{aligned} \frac{dX_1(t)}{dt} &= \beta_{11} X_1(t) + \beta_{12} X_2(t), \\ \frac{dX_2(t)}{dt} &= \beta_{21} X_1(t) + \beta_{22} X_2(t), \end{aligned}\right\} \qquad (2.8)$$

Where

$$\beta_{1J} = \left(\frac{\partial b_1}{\partial x_J} - \frac{\partial m_1}{\partial x_j}\right) \text{ evaluated at } (\alpha, \beta) \qquad (i, j = 1,2).$$

(2.9)

The steady state $(\alpha, \beta)$ of (2.1) is locally asymptotically stable if and only if the trivial solution (0, 0) of (2.8) is asymptotically stable. We can derive the following result from elementary considerations.

In the system (2.1)-(2.6) assume the following:

$$\left.\begin{aligned} \frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1}, \\ \frac{\partial m_1}{\partial x_1} &> \frac{\partial b_1}{\partial x_1} + \frac{\partial m_2}{\partial x_1}, \end{aligned}\right\} \qquad \text{at } (\alpha, \beta)$$

(2.10)

Then the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (2.1) – (2.6).

The proof follows easily if we note that the characteristic equation associated with the linear system (2.8) is given by

Det $\begin{pmatrix} \lambda-\beta_{11} & -\beta_{12} \\ -\beta_{21} & \lambda-\beta_{22} \end{pmatrix} = 0$

Or equivalently,

$$\lambda^2 - \lambda\,(\beta_{11} + \beta_{22}) + \beta_{11}\beta_{22} - \beta_{12}\beta_{21} = 0. \qquad (2.11)$$

It follows from (2.2) and (2.10) that,

$$\beta_{11} < 0, \beta_{11}\beta_{22} - \beta_{12}\beta_{21} > 0$$

And hence the roots of (2.11) have negative real parts implying the asymptotic stability of the trivial solution of (2.8)

If apply the condition (2.10) to the Volterra-Lotka model (2.7) we note that when a nontrivial steady state say $(\alpha, \beta)$ exists for (2.7), (2.10) leads to

$$a_{11}\alpha > a_{21}\,\beta \; and \; a_{22}\,\beta > a_{12}\,\alpha,$$

Which together imply $a_{11}\,a_{22} > a_{12}a_{21}$. If $(\alpha, \beta)$ exists it is known that $a_{11}\,a_{22} > a_{12}a_{21} > 0$ is a sufficient condition for the (local) asymptotic stability of $(\alpha, \beta)$.

The condition (2.10) which guarantees the coexistence of competing populations in (2.1)-(2.6) simply means that the intraspecific self – regulating feedback effects $\partial m_i / \partial x_1\ (i = 1, 2)$ bare higher than its own positive feedback $\partial b_1 / \partial x_1\ (i = 1, 2)$ together with its negative (competitive) effects on its competitor $\partial m_I / \partial x_J\ (i,\ j = 1, 2; i \neq j)$.

## 16. Coexistence under delayed production and interactions, transactions

We will now establish that two competing species with a sufficiently strong intraspecific regulation can coexist in the presence of delayed reproduction and delayed interspecific interactions. It is well known that if there are delays in the self – regulating intraspecific effects, stationary coexistence may not always be possible (Gopalsamy [16], Gopalsamy and aggarwala [15], Shibata and Saito [28]). Delayed reproduction is more general and widespread in most biological populations. To build a competition model with such delays we proceed as follows.

Let $\tau_{ij}\ (i, j = 1, 2)$ be a set of nonnegative constants with $\tau = \max \tau_{ij}\ , i, j = 1, 2,$ and suppose that the two ompeting species display delayed reproduction and delayed interspecific interactions with no delays in the self – regulating negative feedback effects. Such a competition system in a constant environment can be modeled by the autonomous delay differential system

$$\frac{dx_1(t)}{dt} = b_1\big(x_1(t-\tau_{11})\big) - m_1\big(x_1(t), x_2(t-\tau_{12})\big),$$

$$\frac{dx_2(t)}{dt} = b_2\big(x_2(t-\tau_{22})\big) - m\big(x_1(t-\tau_{21}), x_2(t)\big), \qquad (3.1)$$

Where both the creation and destruction rates and the morality rates $b_1$ and $m_1\ (i = 1, 2)$ satisfy the same conditions (2.2)-(2.6) as in the previous section. Along with (3.1) we will

Suppose that the initial population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) sizes are provided in the form

$$x_i(s) = \phi_i(s) \ \text{⓪}, \qquad s \in [-\tau, 0], \qquad i = 1,2. \tag{3.2}$$

(Note that the specification of $\phi_i$ need not be on the entire interval$[-\tau, 0]$; but such a hypothesis is convenient.) Since (3.1) is not of the standard Kolmogorov-type, we have to establish that the solutions of (3.1) will remain nonnegative for nonnegative initial conditions.

To prove the nonnegativity of solutions of (3.1)-(3.2), let $\left\{\psi_1^{(n)}(s)\right\}, s \in [-\tau, 0], i = 1,2; (n = 1,2, \ldots)$ be a sequence of strictly positive continuous functions such that

$$\lim_{n\to\infty} \psi_1^{(n)}(s), = \psi_1(s), \qquad s \in [-\tau, 0], \qquad i = 1,2 \tag{3.3}$$

Let $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ be the solution of (3.1) corresponding to the initial condition

$$\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \left\{\psi_1^{(n)}(s), \psi_2^{(n)}(s)\right\}, \quad s \in [-\tau, 0],. \tag{3.4}$$

Consider the solution $\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\}$ for $t \in [0, \tau]$ (assuming it exists on the whole of $[0, \tau]$). Suppose $x_1^{(n)}$ does not remain nonnegative for all $[0, \tau^*]$; then there exists a $t^* = [0, \tau^*](\tau^* =$ the positive minimum of $\tau_{ij}\ (i, j = 1, 2))$ for which

$$x_1^{(n)}(t^*) = 0 \ \text{and} \ x_1^{(n)}(t)(t^* - \tau_{11})) > 0.$$

It will follow from (3.1) and the positivity of the initial condition that

$$\frac{dx_1^{(n)}(t^*)}{dt} = b_2\left(x_1^{(n)}(t^* - \tau_{11})\right) > 0 \tag{3.5}$$

Where we have used the properties of $b_i$ and $m_i$ from (2.2)-(2.6); it is found that (3.5) contradicts the definition of $t^*$. Thus $x_1^{(n)}(t) > 0$ on $[0, \tau^*]$; similarly $x_2^{(n)}(t) > 0\ on\ [0, \tau^*]$. We can repeat this procedure for intervals of the form $[\tau^*, 2\tau^*], [2\tau^*, 3\tau^*]\ etc$. Thus it will follow that so long as $(x_1^{(n)}(t), x_2^{(n)}(t))$ exists, we have

$$x_1^{(n)}(t) \geq 0\ x_2^{(n)}(t) \geq 0.$$

If we now consider the limits as $n \to \infty$, we get

$$\lim_{n\to\infty}\left\{x_1^{(n)}(t), x_2^{(n)}(t)\right\} = \{x_1(t), x_2(t)\} \tag{3.2}$$

With $x_1(t)$ ⓪,, $x_2(t)$ ⓪ where $(x_1(t), x_2(t))$ is the solution of (3.1)-(3.2) and this is a consequence of the continuous dependence of solutions of (3.1) on initial conditions (Hale [17], page 41).

A second question for (3.1) with (2.2)-(2.6) is whether solutions of (3.1) exist for all $t > 0$? suppose now that solutions of (3.1) do not exist for all $t$ ; then there exists a $t_1 > 0$ such that at least for one of either of $x_1$ *or* $x_2$ we have

$$\lim_{t \to t_1 -} x_i\ (t) = \infty. \tag{3.7}$$

To be specific let us suppose that $\lim_{t \to t_1 -} x_i\ (t) = \infty$; let $t_2$ be the first instant of time for which $x_1(t_2) = \delta_1, t_2 < t_1$. It will follow from (3.1) that

$$\frac{dx_1(t_2)}{dt} = b_1\big(x_1\ (t_2 - \tau_{11})\big) - m_1\big(x_1(t_2), x_2\ (t_2 - \tau_{12})\big)$$

$$< b_1\big(x_1\ (t_2)\big) - m_1\big(x_1(t_2), x_2\ (t_2 - \tau_{12})\big)$$

$$< b_1(\delta_1) - m_1\big(\delta_1, x\ (t_2 - \tau_{12})\big) < 0 \tag{3.8}$$

Which contradicts the definition of $t_2$. Thus solutions of (3.1), (2.2) – (2.6) exist for all $t$ .

We are now ready to consider the effects of time delays in reproduction and interspecific interactions on the coexistence of the two competing species. The steady state $(\alpha, \beta)$ of (2.1) is again a steady state of (3.1). To examine the (local) asymptotic stability of $(\alpha, \beta)$ in (3.1) we let

$$x_1(t) = \alpha + Y_1(t), \quad x_2(t) = \beta + Y_2(t)$$

In (3.1) and derive that the linear variational system governing the perturbations $Y_1, Y_2$ is given by

$$\frac{dY_1\ (t)}{dt} = -m_{11}Y_1(t) + b_{11}Y_1\ (t - \tau_{11}) - m_{12}Y_2\ (t - \tau_{12}),$$

$$\frac{dY_2\ (t)}{dt} = -m_{22}Y_2(t) + b_{21}Y_1\ (t - \tau_{21}) - m_{22}Y_2\ (t - \tau_{22}),$$

Where

$$m_{ij} = \frac{\partial m_1}{\partial x_j}, \quad b_{ii} = \frac{\partial b_i}{\partial x_i}\ at\ (\alpha, \beta). \tag{3.10}$$

It is well known (Bellman and Cooke [3], page 336) that the steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) if the trivial solution of the variational system (3.9) is asymptotically stable. We formulate our first result in the following form whose proof is given in the appendix.

*THEOREM 3.1. Assume that in the two systemic types competition model system (3.1) with delayed density dependent birth and death rates $b_i$ and $m_i$ $(i = 1, 2)$ satisfy the conditions (2.2)-(2.6). Suppose that the conditions (2.10) hold. Then the nontrivial steady state $(\alpha, \beta)$ is (locally) asymptotically stable for (3.1) whatever the delays be $\tau_{ij}$ $(i = 1, 2)$.*

Let $y_1(t) and\ y_2(t)$ respectively denote the population sizes (or biomasses) of a prey species and a predator species which feeds exclusively on the single prey species in a temporally constant environment. In the format of (1.5) the following system will describe the rates of change of the two population sizes:

$$\frac{dy_1(t)}{dt} = f_1(y_1(t)) - g_1(y_1(t), y_2(t)),$$
$$\frac{dy_2(t)}{dt} = f_2(y_1(t)) - y_2(t) - g_2(y_2(t)), \quad (4.1)$$

Where $f_i\ and\ g_i\ i = 1,2)$ respectively denote the birth and death rates of the two species; in order to make (4.1) denote the dynamics of a prey – predator species we will assume that $f_i\ and\ g_i$ satisfy the following conditions (see for instance Rescigno and Richardson [26]); $f_i\ and\ g_i$ are continuous with continuous partial derivatives for all $y_i \geq 0\ (i = 1, 2); also\ we\ require$

$$\frac{\partial f_1}{\partial y_1} > 0; \quad \frac{\partial g_1}{\partial y_1} \geq 0; \quad \frac{\partial g_1}{\partial y_2} \geq 0,$$
$$\frac{\partial f_2}{\partial y_1} \geq 0; \quad \frac{\partial g_2}{\partial y_2} \geq 0; \quad \frac{\partial g_2}{\partial y_2} \geq 0, \quad \text{Where } y_1 \geq 0, y_1 \geq 0, \quad (4.2)$$

(xxix)

$$f_1(0) = 0, \quad g_1(0, y_2 \equiv 0,$$
$$f_2(0, y_2) = 0 = f_2(y_1, 0); \quad g_2(0) = 0, \quad (4.3)$$

(xxx) For some constant $y_1^* > 0$ we have

$$f_1(y_1^*) - g_1(y_1^* 0) = 0, \quad (4.4)$$

(xxxi) There exists a point $(\alpha^*\beta^*), \alpha^* > 0, \beta^* > 0$ such that

$$f_1(\alpha^*) - g_1(\alpha^*, \beta^*) = 0,$$
$$f_2(\alpha^*, \beta^*) - g_2(\beta^*) = 0, \quad (4.5)$$

(xxxii) There exist positive constants $n_1, n_2$ such that

$$f_1(\eta_1) - g_1(\eta_1, y_2) < 0, y_2 \geq 0,$$
$$f_2(\eta_1, \eta_2) - g_2(\eta_2) < 0 \quad (4.6)$$

The conditions (4.2) mean that the system (4.1) is of the oppugnacy systems type; (4.3) will imply that (0,0) is a trivial steady state of (4.1); (4.4) means that in the absence of the predator, the prey species has a nontrivial (positive) steady state;

(4.5) implies that (4.1) has a nontrivial steady state in the interior of the positive quadrant of the state species of (4.1); (4.6) can be shown (see below) to imply that no credit type and debit type system types can grow unboundedly.

An example of (4.1)-(4.6) is the familiar Volterra-Lotka prey-predator system

$$\frac{dy_l(t)}{dt} = r_l y_l(t) - [a_{ll} y_l^2(t) + a_{12} y_l(t) y_2(t)],$$
$$\frac{dy_2(t)}{dt} = a_{12} y_1(t) y_2(t) - [r_2 y_2(t) + a_{22} y_2^2(t)],$$

Where $r_{l,}\ a_{ij}\ (i,j=1,2)$are positive costants such that

$$a_{21}r_1 - r_2a_{11} > 0.$$

As in the previous case of competition we have to verify that when $y_1(0) > 0, y_2(0) > 0$ the corresponding solution of (4.1) will be nonnegative for $t \geq 0$ for which the solution exist. Since the co-ordinate axes $y_1 = 0, y_2 = 0$ are invariant sets for (401) by virtue of (4.3), it will follow that the solution of (4.1) starting from the interior of the positive quadrant cannot reach the outside of that quadrant. This is a consequence of the invariance of the co-ordinate axes for (4.1)

To prove that solutions of (4.1)-(4.6) exist for all t $\geq 0$ let us suppose that $y_l 0\ and\ y_2(0) satisfy\ 0 < y_l\ (0)$ satisfy $0< y_l(0) < \eta_l, 0 < y_2(0) < \eta_2$ ( on the other hand if either one of $y_l(0) \geq \eta_l and\ y_2(0) \geq \eta_l$ or both hold, it will follow that the corresponding $y_l(t)$ will decrese by virtue of (4.6) at least for small $t \geq 0,$ ). Suppose $y_l(t)$ does not exist for all $t \geq 0$; then for some $t_l < \infty\ lim_{t\to t_l-} y_l(t) = \infty$and let $t_2 > 0$ be the first instant for which $y_1(t_2) = \eta_1$; then we have from (4.1) and (4.6) that

$$\frac{dy_1(t_2)}{dt} = f_1(\eta_1) - g_1(\eta_{1,}\ y_2) < 0\ (assumimg\ y_1(0) < \eta_1)$$

Which contradicts the definition of $t_2$ showing that $y_1(t)$ exists for all $t \geq 0$; if $y_2(t)$ is not defined for all $t \geq 0$, then for some $t_3 \geq \infty$ we will have $lim_{t-t_3-} y_2(t) = \infty$;now let $t_4$ be the first instant for which $y_2(t_4) = \eta_2$; then from the properties of $f_2$ we have

$$\frac{dy_2(t_4)}{dt} = f_2(y_1(t_4), \eta_2) - g_2(\eta_2) \leq f_2(\eta_1, \eta_2) - g_2(\eta_2)$$

$$< 0 (by\ (4.6) assuming\ 0 < y_1(0) < \eta_1, 0 < y_2(0) < \eta_2)$$

Which again contradicts the definition of $t_4$. Thus $y_2(t)$ is defined for all $t \geq 0$.

We will now derive a set of sufficient condition for the (local) asymptotic stability of the steady state $(\alpha^*, \beta^*)$ of (4.1). in the next section we will show that under the same conditions and in the presence of various delays , the steady state $(\alpha^*, \beta^*)$remains (locally) asymptotically stable.

To examine the asymptotic stability of the nontrivial steady state $(\alpha^*, \beta^*)$ in (4.1) we let

$$y_1(t) = \alpha^* + U_1(t), \qquad y_2(t) = \beta^* + U_2(t)$$

In (4.1) and derive the liner variational system in the perturbations $U_1, U_2$ as follows:

$$\left.\begin{aligned} \frac{dU_1(t)}{dt} &= \frac{\partial f_1}{\partial y_1}U_1(t) - \frac{\partial g_1}{\partial y_1}U_1(t) - \frac{\partial g_1}{\partial y_2}U_2(t), \\ \frac{dU_2(t)}{dt} &= \frac{\partial f_2}{\partial y_1}U_1(t) - \frac{\partial f_2}{\partial y_2}U_2(t) - \frac{\partial g_2}{\partial y_2}U_2(t), \end{aligned}\right\} \quad (4.7)$$

Where all the partial derivatives in (4.7) are evaluated at $(\alpha^*, \beta^*)$. The steady state $(\alpha^*, \beta^*)$ of (4.1) is (locally) asymptotically stable if and only if the trivial solution of (4.7) is asymptotically stable which will be the case if and only if the roots of the characteristic equation

$$\text{Det} \begin{pmatrix} \lambda - (f_{11} - g_{11}) & g_{12} \\ -f_{21} & \lambda - (f_{22} - g_{22}) \end{pmatrix} = 0,$$

Where $f_{ij} = \frac{\partial f_1}{\partial y_l}$, $g_{ij} = \frac{\partial g_1}{\partial y_j}$, at $(\alpha^*, \beta^*)$ (4.8)

$(i,j = 1, 2)$

Have negative real parts. If $\lambda_1, \lambda_2$ are the riits of (4.8) then

$$\lambda_1 + \lambda_2 = (f_{11} - g_{11}) + (f_{22} - g_{22}), \quad (4.9)$$

$$\lambda_1 \lambda_2 = (f_{11} - g_{11})(f_{22} - g_{22}) + f_{21} g_{12}.$$

The proof of the following result is an immediate consequence of (4.7)-(4.9)

THEOREM 4.1 *In the CREDIT-DEBIT system (4.1)-(4.6) assumes that the partial derivatives satisfy the following additional conditions:*

$$\frac{\partial g_1}{\partial y_1} > \frac{\partial f_1}{\partial y_1} + \frac{\partial f_2}{\partial y_1}, \quad \frac{\partial g_2}{\partial y_2} > \frac{\partial f_2}{\partial y_2} + \frac{\partial g_1}{\partial y_2}, \quad \text{Evaluated at } (\alpha^*, \beta^*), \quad (4.10)$$

*Then the steady state $(\alpha^*, \beta^*)$, is (locally) asymptotically stable for the prey- predator system.*

## 12. Delays in reproduction, predation and gestation

A number of authors have proposed and analysed several models of prey- predator systems with time delays. Most of the models have been derived from the so called models of the Kolmogorov-type f (1.2) with time delays incorporated in the average growth rates of the prey and predator species. Volterra [32] proposed the following model for a prey-predator system. The equations are used for the formulation of Credit-debit system of atoms, accounts, economic variable, cosmological variable, astrophysical; variables, and qubit systems for which the model is being applied.

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \delta_1 N_1(t) - \int_{\infty}^{t} F_1(t-s) N_2(s) ds],$$

(5.1)

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 + \delta_2 N_1(t) - \int_{\infty}^{t} F_2(t-s)N_1(s)ds],$$

Where $r_1, r_2, \delta_1, \delta_2$ *are positive constant and* $F_1, F_2$ is nonnegative continuous delay kernel defined and integrable on $[0, \infty]$. (5.1) has no predator self-regulating negative feedback. Brelot [6] considered a modified format of (5.1) in the form

$$\frac{dN_1(t)}{dt} = N_1(t)[r_1 - \lambda_1 N_1(t) - \delta_1 N_2(t) - \int_{-\infty}^{t} F_1(t-s)N_2(s)ds],$$

$$\frac{dN_2(t)}{dt} = N_2(t)[-r_2 - \lambda_2 N_2(t) - \delta_2 N_1(t) - \int_{-\infty}^{t} F_2(t-s)N_1(s)ds], \quad (5.2)$$

Which has term representing intraspecific competition among both predator and prey species. Under suitable condition on the various parameters in (5.2) one can show that (5.2) has a (locally) asymptotically stable steady state. A number of intergrodifferential equation models of the types (5.1) and (5.2) have been investigated by Cushing [12]

Starting from Hutchinson's (1948)[19] delay logistic equation, May[23] has proposed the following system

$$\frac{dN_1(t)}{dt} = rN_1(t)\left[1 - \frac{N_1(t-\tau)}{K}\right] - \alpha N_1(t)N_2(t),$$

$$\frac{dN_2(t)}{dt} = -bN_2(t) + \beta N_2(t)N_1(t), \quad (5.3)$$

Where r, $\tau$, K, $\alpha, \beta, b$ are positive constants; (5.3) contains a single discrete delay; one can modify (5.3) and incorporate a continuously distributed delay in it so that (5.3) can be written as (May [23])

$$\frac{dN_1(t)}{dt} = rN_1(t)[1 - \int_{-\infty}^{t} F(t-s)N_1(s)ds] - \alpha N_1(t)N_2(t),$$

$$\frac{dN_2(t)}{dt} = -bN_2(t) + \beta N_1(t)N_2(t). \quad (5.4)$$

The first model of a prey-predator type which departs from the Kolmogorov-type formulation is due to Wangersky and Cunningham [33] who considered the model system

$$\frac{dN_1(t)}{dt} = a_1N_1(t) - b_1N_1^2(t) - c_1N_1(t) - N_2(t),$$

$$\frac{dN_2(t)}{dt} = a_2N_2(t) + c_2N_1(t-\tau)N_2(t-\tau), \quad (5.5)$$

Where $a_1, a_2, b_1, c_1, c_2, \tau$ are positive constants; (5.5) means that a duration of $\tau$ time units elapses when an individual prey is killed and the moment when the corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) Recently Cushing [13] has considered a model of the form

$$\frac{dN_1(t)}{dt} = rN_1(t)\left[1 - \frac{N_1(t)}{k} - \alpha N_2(t)\right],$$

$$\frac{dN_2(t)}{dt} = -d_2^*(t) + bN_1(t)\int_0^\infty \delta(a)N_2(t-a)e^{-d*a}\,da,$$

(5.6)

Where $r, k, \alpha, d^*, b$ are positive constants and $\delta$ is related to the age dependent fecundity of the predator species.

With the models (5.1)-(5.6) in the background let us now consider a time delayed model similar to (4.1) in the following form:

$$\frac{dN_1(t)}{dt} = f_1\big(N_1(t-\tau_{11})\big) - g_1\big(N_1(t), N_2(t-\tau_{12})\big),$$

$$\frac{dN_2(t)}{dt} = f_2\big(N_1(t-\tau_{21})\big), N_2(t-\tau_{12}) - g_1(N_2(t)),$$

(5.7)

Where $\tau_{ij}$ $(i,j=1,2)$ are nonnegative constants with $\tau = \max \tau_{ij}, i,j = 1,2$; the birth and death rate functions $f_1, f_2, g_1, g_2$ satisfy the same conditions as in (4.1)-(4.6). Along with (5.7) we specify the initial conditions as follows:

$$N_1(s) = \phi_1(s) > 0,$$

$$s \in [-\tau, 0].$$

$$N_1(s) = \phi_1(s) > 0,$$

We3 will assume that in addition to (4.2) we have for (5.7) that

$$\frac{\partial f_2}{\partial N_1} > 0, \frac{\partial f_2}{\partial N_2} > 0, f_2(N_2, N_2) > 0 \; whenever \; N_1 > 0, \; N_2 > 0. \; (5.9)$$

It is not difficult to show by the methods previously used for the competition system that whenever $N_1(s) > 0$ and $N_2(s) > 0$ for $s \in [-\tau, 0]$, solution of (5.7) exist for all $t \geq 0$ and remain nonnegative for all $t \geq 0$. also $(\alpha^*, \beta^*)$ of (4.5) is a steady state of (5.7).

We will discuss the form and meaning of the delays in (5.7) in the next section. The effects of the various delays in (5.7) on the (local) asymptotic stability of $(\alpha^*, \beta^*)$ can be derived from the following theorem whose proof is quite similar to that of Theorem 3.1 and hence the details of the proof are omitted.

THEOREM 5.1 *In the time delayed prey-predator system(* corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) *) (5.7) let the conditions (4.2)-(4.6), (5.8) and (5.9) hold for the birth and death rates; furthermore suppose the condition (4.10) hold. Then for all nonnegative delays in (5.7), its nontrivial steady state $(\alpha^*, \beta^*)$ is (locally) asymptotically stable.*

## 6. Estimation of decay rates

We will first derive a convention sufficient condition for the asymptotic stability of the trivial solution of a linear system of the form

$$\left.\begin{aligned} \frac{dY_1(t)}{dt} &= -m_{11} Y_1(t) + b_{11} Y_1(t-\tau_{11}) - m_{12} Y_2(t-\tau_{12}), \\ \frac{dY_2(t)}{dt} &= -m_{22} Y_1(t) + m_{21} Y_1(t-\tau_{21}) - b_{22} Y_2(t-\tau_{22}), \end{aligned}\right\} \; t > 0, \quad (6.1)$$

Where $m_{1J}, b_{i1}$ $(i, j = 1,2)$ are positive constants while $\tau_{1J}(i, j = 1,2)$ are non negative with $\tau = \max(\tau_{1J})(i, j = 1,2$.as is shown in the appendix, the following system dominates (6.1*)*

$$\left.\begin{aligned} \frac{dZ_1(t)}{dt} &= -m_{11} Z_1(t) + b_{11} Z_1(t-\tau_{11}) - m_{12} Z_2(t-\tau_{12}), \\ \frac{dZ_2(t)}{dt} &= -m_{22} Z_1(t) + m_{21} Z_1(t-\tau_{21}) - b_{22} Z_2(t-\tau_{22}), \end{aligned}\right\} \; t > 0, \quad (6.2)$$

$Z_i(s) = |Y_i(s)|, \qquad i = 1,2 \;\; for \; s \leq 0 \;\; for \; which \; Y_i(s) \; is \; defined.$

Thus the asymptotic stability of the trivial solution of (6.2) implies that of (6.1). it is also shown in the appendix that the solutions of (6.2) are nonnegative for all t $\geq 0$. we will use the following vector-matrix notation:

$$Z = (Z_1, Z_2)^{\tau}, \qquad ||Z|| = |Z_1(t)| + |Z_2(t)| = Z_1(t) + Z_2(t)$$

$$B = \begin{bmatrix} b_{11} & m_{12} \\ m_{21} & b_{12} \end{bmatrix}, \quad ||B|| = max \begin{matrix} b_{11} + & m_{21} \\ b_{21} + & m_{12} \end{matrix}, \tag{6.3}$$

$$A = \begin{bmatrix} -m_{11} & 0 \\ 0 & -m_{22} \end{bmatrix}.$$

The logarithmic norm or the measure of any square matrix B, denoted by $\mu(B)$is defined as follows:

$$\mu(B) = \lim_{\theta \to 0+} \frac{/\!/I + \theta B/\!/ - 1}{\theta}$$

($I$ being an identity matrix of the same size as $B$).The matrix measure compatible with the norms in (6.3) is given by

$\mu(B) = \max_J [b_{jj} + \sum_{\substack{i=1 \\ i \neq J}}^{2} | b_{jJ} | ]$ $B=(b_{1J})$.

We refer to Coppel [11]or Vidyasager [31]for a number of properties of matrix measure We have shown in the appendix that the derivative of $H(Z_1,Z_2)(t)$(seeA.10)satisfies the relation

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq [-m_{11}+b_{11}+m_{21}]Z_1(t)+[-m_{22}+b_{22}+m_{12}]Z_2(t)$, and hence we have

$\frac{d}{dt} H(Z_1,Z_2)(t) \leq \{\mu(A) + /\!/B/\!/\}\{ Z_1(t)+ Z_2(t)\}$

If we assume for the matrix norm $\|B\|$ as in (6.3)that

$$\mu(A) + /\!/B/\!/ < 0 \; or \; -\mu(A) > /\!/B/\!/$$

Then from the results of the appendix (Lemmas 5 and 6) it will follow that $Z_1(t)+Z_2(t) \to 0$ as $t \to \infty$ and thus we have the following:

Lemma 6.1 Assume that the Coefficients of the linear system (6.1)satisfy the condition (6.7). Then the trivial solution of(6.1)is asymptotically stable.

When (6.7)holds an estimate for the decay rate of solutions of (6.1)can be obtained as follows A necessary and sufficient condition for the asymptotic stability of the trivial solution is that all the roots of the characteristic equation

Det $G(\lambda,T_{iJ})=0$

Have negative real parts where

$G(\lambda,T_{1J})=[\begin{matrix} \frac{\lambda+m_{11}}{0} & \frac{0}{\lambda+22} \end{matrix}] - \begin{matrix} b_{11}{}^{e-\lambda T}{}_{11} & -m_{12}{}^{e-\lambda T}{}_{12} \\ -m_{21}{}^{e-\lambda T}{}_{21} & b_{22}{}^{e-\lambda T}{}_{22} \end{matrix}$

$=\lambda I - A - B(\lambda, T_{ij})$(say),

A positive number $\mathfrak{a}$(if it exists)is said to be a decay rate of solution of (6.1) if and only if every root of (6.8)satisfies the condition

$\mathrm{Re}(\lambda)+\mathfrak{a} < 0.$

When (6.11)holds the solution of (6.1)will be of the order of $e^{-at}$ and hence the supremum of all such $\mathfrak{a} > 0$ $a$ satisfying (6.11)will be a desirable parameter to estimate However such a supremum of all such (6.11)may not be possible to determine since it has to be done directly from (6.8)-(6,10)(see also Section 3).We will obtain an estimate of $\mathfrak{a}$ (not necessarily the supremum)I (6.11)by the following procedure.

Replace $\lambda$ byz-$\mathfrak{a}$ in (6.8)and define

$G(z-\mathfrak{a}, T_{IJ}) = M(z) = (z-a)I - A - B(z-\mathfrak{a}, T_{IJ})$.

By Lemma 6.1 a Sufficient condition for all roots of M(z)=0 to have negative real parts equivalently to satisfy

$\mathrm{Re}(\lambda) < -a$

is the following :

$-\mu(aI + A) > \|B_1(a, T_{ij})\|,$

Where

$$B_1(a, T_{ij}) = \begin{pmatrix} b_{11}^{\,e-\lambda T_{11}} & -m_{12}^{\,e-\lambda T_{12}} \\ -m_{21}^{\,e-\lambda T_{21}} & b_{22}^{\,e-\lambda T_{22}} \end{pmatrix}$$

Since $\mu(\mathfrak{a} I + A) = \mathfrak{a} + \mu(A)$ we have from (6.13) that

$\mathfrak{a} + \mu(A) + \|B_I(\mathfrak{a}, T_{IJ})\| < 0$

Is a sufficient condition for Re $(\lambda) < -a$If we denote the left side of (6.15)by $\mathfrak{a} f(\mathfrak{a}, T_{Ij})$ we note that

$f(0, T_{Ij}) = \|B\| + \mu(A) < 0$ $if$ (6.7)$holds$;

Al so $\partial f / \partial \mathfrak{a} > 0, f$ is continuous and monotonic in $(\mathfrak{a})$;$\to \infty as\ \mathfrak{a} \to \infty$; thus by the continuity of $f$ on $\mathfrak{a} \geq 0$, there exists an interval of the from $[0, ^{*})$such that

$f(0, T_{Ij}) < 0$ $for\ \mathfrak{a} \in [0, \mathfrak{a}^{*})$

And hence

$f(\mathfrak{a}, T_{Ij}) = 0$

Thus an upper bound $\mathfrak{a}^*$ of $\mathfrak{a}$ for which we have $f(\mathfrak{a}, T_{Ij}) < 0$ $\mathfrak{a} \in [0, \mathfrak{a}^*)$ can be obtained as the unique (by the monotonicity of $f$ on $\mathfrak{a}$) positive root ofmax $b_{11}^{e-\lambda T}{}_{11}$ - $m_{12}^{e-\lambda T}{}_{21}$ + $\mathfrak{a}$-min $m_{11}$ =0

$b_{22}^{e-\lambda T}{}_{22}$ - $m_{21}^{e-\lambda T}{}_{12}$ $m_{22}$

Where the maximum in (6.18) is taken over $b_{11}, b_{22}, m_{12}, m_{21}$ and $T_{ij}$ $(i, = j\ 1,2)$. A numerical procedure can be devised to solve equations of the form (6.18) in concrete cases. If one is interested in a smaller bound on the decay rate it can be obtained as the unique root of

$\| B \| e^{ar} + a + \mu(A) = 0, \tau = \max_{i,J} T_{ij}$. (6.19)

It will follow from (6.19) than an increase in T

Will reduce the value of the root of (6.19). We remark that since our estimate $a^*$ of dacay rate is based on a sufficient condition of the form (6.7), we can conclude that the solotions of (6.1) will converge at least as fast as $e^{-a*t}$ where $a^*$ is the unique root of either of (6.18) or (6.19).

## Implications

Authors have formulated a class of two systems model systems staring from the elementary principal that the rate of change of a population confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) equals the creation rate minus the destruction rate. This is a departure from the conventional model building in population ecology where one usually starts from the Malthusian formulation, now known as the Kolmogorov formulation in the nonlinear context. In our analysis of the models proposed we have obtained sufficient conditions for the global existence of nonnegative solutions of the model systems (this is trivial in the Kolmogorov formulation) and the other desirable properties such as the existence of suitable steady states. The two classes of models proposed for competition and debt contain intraspecific self – regulation biologists that each systems is under the action of a self-regulating negative feedback force, opinions and arguments differ regarding the mechanism by which such an effect is realized; we refer for more details on this aspect to a collection of articles compiled by Tamarin [30].

Gopalasamy's main result is the following : if a self –regulating negative feedback mechanism (intraspecific oppugn ant) can act without time delay and if such an effect can keep an ecosystem ( a community of one or more species of populations) asymptotically stable ( at least locally about a non trivial steady state, then time delays of any magnitude in any other interaction involving positive or negative feedbacks cannot destabilize such a stable ecosystem ( in a mathematical sense); that is such delays cannot cause a hopf-type bifurcation leading to persistent and undamped fluctuations. The delay-induced bifurcations to oscillations in models have been investigated by many authors. (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Qualitatively result is not new to the literature the effects of delays on otherwise stable and unstable systems have been considered by Cushing[12]where (pages 35-39)he has shown that an otherwise unstable system can be made

less unstable by delays, Such an effect of delays on unstable systems appear to have been first proposed by Beddinton and May [2]who have also shows that delays can make stable systems less stable leading to a conclusion "time delays are not necessarily destabilizing" It is also shown by Cushing([12], pages 75-103)that while time delays cannot stabilize(in the mathematical sense)an otherwise unstable two species system in competition or mutualism time delays in there species systems can stabilize otherwise unstable systems (even in a mathematical sense)Cushing and Saleem [14]have formulated a model where delays in the positive feedback can have a stabilizing effect: Hale[18]discusses briefly the stabilizing nature of delays and given an example to show that an increase in delays can make an unstable system stable Thus for a better understanding of the effects of delays on otherwise stable systems one has to estimate the rate of decay of perturbations (or the characteristic return time )in terms of the parameters of the system especially when the delays do not cause a mathematical instability This aspect has been investigated by a procedure slightly different from that of the characteristic equation we have proposed (section6) method for the estimation of decay rate of perturbations when there are multiple delays Since many species systems with multiple delays cannot be analysed (due to complexity)by the characteristic equation our methods using comparison lead to easily verifiable sufficient conditions The sufficient conditions obtained are in the form of a diagonal dominance type; these conditions can be weakened to quasi-diagonal dominance type as in Ladde [20], Siljak[20]Anderson [1], Lewis and Anderson [21,22](Ladde's result [20] is delay dependent due to the fact that delay appears in self –regulation also).Verification of quasi –diagonal dominance conditions in arbitrary matrices (except for Metzler matrices )is an impossible task and quasi –diagonal dominance is hard to interpret.

If follows from our discussion that those confederation systems which display delay induced instability leading to permanent and undamped oscillations should have one or more of the following aspects:

(1)time delays are in intraspecific interaction terms (and possibly in other terms also): (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator –population.)

(2) Intraspecific self-regulation terms are absent in density dependent from in one or more populations of the system;

(3) Intraspecific negative feedback effect are not strong enough to maintain the stability of the relevant steady state of the nondelayed system.

We can now infer how the delay induced stability leads to fluctuations; a primary mechanism of such fluctuations can be described at least theoretically as follows : in its attraction of a system to a steady state, intraspecific self-regulation fails to be effective and the system overhoots the steady state and at such higher densities intraspecific competition becomes operationally effective leading the system to understood the steady state and this process is repeated again and again resulting in undamped oscillations about the steady state. It follows from our results that if intraspecific competition is sufficiently strong, the system cannot become oscillatory even when there are time delays, so long as they do not interfere with the strength of the intraspecific self – regulation.

We are thus led to the question of existence and importance of intraspecific competition in population systems. In the introduction to an anthology of articles by several authors, Tamarin [ 30] begins with the statement "all animal populations have the reproductive capacity to increase without limit, yet none di" and in his concluding remarks writes " this leaves us with only one great concept in the area of confederation do main regulation: that all credit-debit are regulated" More specifically this is applicable to all the confederation of (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions,

financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

In his discussion on the effect of competition on the densities of populations, Nicholson ([ 25], $page$ [ 147] ) has reached the conclution "Intraspecific competition automatically regulates the severity of its action to the requirements of each denizen , provided the inherent resistance of the environment is sufficiently low to permit the species to exist." In order to look for a mechanism by which intraspecific competition plays its role of decreasing growth rate as confederation density increases, Calhoun [ 7] set up an experiment with Norway rates (Rattus norvegicus) and found that increased density leads to increased hostility which disturbs the psychology and physiology causing decreased reproductive output. Many other mechanisms by which intraspecific compition self-regulate populations have been put forward by a number of authors (see Tamarin [30]). This is true of (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

If we accept the view that intraspecific competition is really predominant in population regulation, then we can understand how and why time delays destabilize stable system leading to oscillations; time delays render intraspecific competitive regulation territorially (in case of particles we take the field in which it operates) weak so that self –regulatory forces become temporarily insufficient. If we now agree that all population models should incorporate self-regulating intraspecific competition terms in surfeit strength along with other tropic intraspecific competition terms in sufficient strength along with other atrophic interactions then it is not fair to say " time delays destabilize" and we note from our analysis that it is usually a failure in self- regulation which can cause instability rather than delays by themselves.

We will give one example to illustrate our point. Consider a credit-debit system in which the debits is self –regulating by a logistic load as follows: (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

$$\frac{dx}{dt} = r\left(x - \frac{x^2}{K}\right) - axy = b_I(x) - m_I(x,y) \quad (say),$$

$$\frac{dy}{dt} = \beta xy - \delta y = b_2(x,y) - m_2(y) - \quad (say),$$

Where $K, a, \beta, \delta$ are positive constants such that the steady state $x^*\ y^*$ exists where

$$x^* = \frac{\delta}{\beta}; \quad y^* = \frac{r}{a}\left(-\frac{\delta}{K\beta}\right) > 0.$$

In this example the predator self –regulation $(\partial m_2/\partial y)$ is not sufficiently strong in the sense

$$x^* = \left(\frac{\partial m_2}{\partial y}\right)_{(x^* y^*)} \succ \left(\frac{\partial b_2}{\partial y} + \frac{\partial m_1}{\partial y}\right)_{(x^* y^*)}$$

If in such a model there is a time delay in the credit creation rate, such a delay can lead to instability of $(x^*\ y^*)$ of setting up undamped oscillations. If however we modify the growth rate by the inclusion of a density dependent self-regulation so that credit

$$\frac{dy}{dt} = \beta\, xy - \delta y - \eta y^2$$

For some positive constant η then it is possible, depending on η, for such a modified system to remain stable without oscillations even if there is a time delay in the credit creations rate. We have another completely different justification to have $\eta y^2$ in the credit dynamics: for instance if the debit is removed from this cannot happen. While this example serves to illustrate an aspect of our result, our advocacy of the significance is not based on this example.

Gopalasamy has thus derived conclusion that "time delays need not destabilize" in models with several discrete delays. However we have provided in the appendix the necessary tools (Lemmas 3,4,5) to prove all our results when the relevant delays are all (or some) continuously distributed. Some authors (Caswell [9], May [23], Cushing [12] appear to regard continuously distributed dalays as financially/cosmologically credit wise obtained when he subjected algae to a variable nitrate environment and obtained a better fit for his data with a continuous delay. The questions of whether a discrete or continuous delay is more realistic may well depend on what process is being modeled. For instance an increased supply of some resource at some point in time may trigger reproduction at a particular moment in time later irrespective of the conditions of the resource during the intervening period. If the duration of increased resource supply necessary to elicit such a response is short compared to the intervening period before the reproduction occurs, then a discrete time delay may be appropriate. In a large number of cases of animal. Note fungus algae are prey. Here it shall be an account which loses the energy, amount, or matter for that matter. Same result holds good for (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.)

Populations of particles, atoms, molecules, propositions, systems types chosen the past dietary and nutritional history of an animal over a long period plays an influential role in determining the current bahaviour of the dietary and nutritional history of an animal over a long period plays an influential role in determining the current behavior of the animal and in such cases continuous time delays will be appropriate.

In conclusion the following comments are appropriate. The time delays since the current creation rate is usually due to a contribution of those mature animals living a few units of time earlier depending on the type of animal or organism; the time delays appearing in interspecific interactions can be due to the fact that in order to successfully interact and compete interferentially a developmental maturity may be necessary. In the case of predation, delays in the birth rate of prey can be due to delayed maturation while that in the birth rate of a predator is due to gestation delay with delayed fecundity; the delays in the predation rate are again are again as before due to the maturity needed for a predator to successfully organize and succeed in an attack.

In our analysis we have assumed throughout a temporally uniform environment neglecting dispersal of populations and their age structure. We expect our results to hold under suitable conditions both in temporally periodic environments with the steady state of a constant environment being replaced by a periodic solution

(nonstationary steady state) and when the effects of dispersion and age structures are considered. It is worthwhile to modify our purely ecological, cosmological, financial, propositional, problematic systems models so as to provide them with a genetic foundation and examine the effects of time delays since forces of selection act with some time delay. Such a generalization will definitely lead to a better understanding of "how confederations regulate themselves".

## Appendix

First we derive a comparison result.

Lemma 1. Consider the delay differential system

$$\left.\begin{aligned}\frac{dY_1(t)}{dt} &= -m_{11}\ Y_1(t) + b_{11}Y_1(t-\tau_{11}) - m_{12}\ Y_2(t-\tau_{12}),\\ \frac{dY_2(t)}{dt} &= -m_{22}Y_2(t) - m_{21}Y_1(t-\tau_{21}) + b_{22}\ Y_2(t-\tau_{22}),\end{aligned}\right\} (A.1),$$

Where the $m_{ij}$, $b_{ij}$ $(i.j = 1,2)$ are positive while $\tau_{ij}$ $(i.j = 1,2)$ nonnegative with $\tau = \max \tau_{ij}$, $i,j = 1,2$. Then

$$\left.\begin{aligned}\frac{d|Y_1(t)|}{dt} &\le -m_{11}\,|Y_1(t)| + b_{11}|Y_1(t-\tau_{11})| + m_{12}\ Y_2(t-\tau_{12})|,\\ \frac{d|Y_1(t)|}{dt} &\le -m_{22}\,|Y_2(t)| + m_{21}|Y_1(t-\tau_{21})| + b_{22}\ |Y_2(t-\tau_{21})|,\end{aligned}\right\} (A.2)$$

Where

$$\frac{d}{dt}\,|\,Y_1(t)\,| = \left\{\begin{aligned}&\frac{dY_1(t)}{dt}\ if\ \ Y_1(t) \ge 0,\\ &\frac{dY_1(t)}{dt}\ if\ \ Y_1(t) < 0,\end{aligned}\right\}\ i = 1,2. \qquad (A.3)$$

The proof is by direct verification.

LEMMA 2. Consider now (A.2) and the following system

$$\left.\begin{aligned}\frac{dZ_1(t)}{dt} &= -m_{11}\ Z_1(t) + b_{11}Z_1(t-\tau_{11}) + m_{12}\ Z_2(t-\tau_{12}),\\ \frac{dZ_2(t)}{dt} &= -m_{22}Z_2(t) + m_{21}Z_1(t-\tau_{21}) + b_{22}\ Z_2(t-\tau_{22}),\end{aligned}\right\} \qquad (A.4)$$

Together with the initial conditions

$$Z_1(s) > |Y_1(s)|,\ \ s \in [-\tau, 0],\ i = 1,2. \qquad (A.5)$$

Then

$$|Y_1(t)| < Z_1(t) \text{ for } t \ge 0 \quad \text{and} \quad i = 1,2. \qquad (A.6)$$

PROOF. Let $U_1(t) = Z_i(t) - |Y_1(s)|$, $i = 1,2.$ $t \geq -\tau.$ Then we have from (A.2) and (A.4),

$$\left.\begin{aligned} \frac{dU_1(t)}{dt} &\geq -m_{11}\ U_1(t) + b_{11}U_1(t-\tau_{11}) + m_{12}\ U_2(t-\tau_{12}), \\ \frac{dU_2(t)}{dt} &\geq -m_{22}U_2(t) + m_{21}U_1(t-\tau_{21}) + b_{22}\ U_2(t-\tau_{22}), \end{aligned}\right\} t > 0, \quad \text{(A.7)}$$

With

$$U_1(s) = Z_i(s) - |Y_i(s)| > 0, \quad i = 1,2, s \in [-\tau, 0]. \qquad \text{(A.8)}$$

Consider (A.7)-(A.8) first on the interval $[\tau, 0]$ : It will follow from (A.7) that $U_1(t) \geq U_1(0) \exp\ [-m_{11}t]$

$$+\int_0^t \left[b_{11}U_1(s-\tau_{11}) + m_{12}\ U_2(s-\tau_{12})\right] \exp\left[-m_{11}(t-s)ds > 0,\right] \qquad \text{(A.9)}$$

For all $t \in [0, \tau^*]$ since $U_i(s) > 0$ on $[-\tau^* 0]$ where $\tau^*$ = positive minimum of $(\tau_{11,}\tau_{12,}\tau_{21,}\tau_{22,})$. Similarly $U_2(t) > 0$ for $t \in [0, \tau^*]$. Now by a similar argument it will follow that $U_1(t) > 0, U_2(t) > 0$ for $t \in [\tau^*, 2\tau^*]$. Repeating the above procedure we obtain the result (A.6).

It follows from Lemmas 1 and 2 that in order to prove theorem 3.1 it is enough to prove that $Z_i(t) \to 0$ as $t \to \infty$ for $i = 1,2$ . which we will do as follows.

PROOF OF THEOREM 3.1 Define a Lyapunov-like function

$$H(Z_1,Z_2)(t)=Z_1(t)+Z_2(t)$$

$$+\int_{1-T}^t b_{11}\ Z_1(s)ds + \int_{1-T}^t m_{12}\ Z_2(s)ds$$

$$+\int_{1-T}^t m_{21}\ Z_1(s)ds + \int_{1-T}^t b_{22}\ Z_{22}(s)ds$$

Since each term of H in (A.10)is nonnegative for t>0, if we show that $H \to 0$ as $t \to \infty$ it will follow from Lemma 2 that

$$0 \leq |\ Y_{1,}(t)\ | \leq Z_1(t) \to 0 \text{ as } t \to \infty,$$

Implying that the perturbations $Y_{1,}$, $Y_2$, governed by (A.1)approach zero as $t \to \infty$, and this will prove our result Calculate the rate of change of H along the solutions of (A.4)so that

$$\frac{dH}{dt} = \frac{\partial H}{\partial t} + \frac{\partial H}{\partial Z_1}\frac{dZ_1}{dt} + \frac{\partial H}{\partial Z_2}\frac{dZ_2}{dt}$$

$$=-[\,Z_1(t)\{m_{11} - (b_{11} + m_{12})\} + Z_2(t)\{m_{22} - (b_{22}m_{12})\}] < 0$$

On using (2.10)since $Z_1(t),Z_2(t)$ are positive for all t>0. Thus $H(Z_1, Z_2)(t) \to 0$ as $t \to \infty$ since H is bounded below by zero: it follows from Lemma 6 below that $Z_1(t) \to 0, Z_2(t) \to 0$ as $t \to \infty$ the proof is complete.

The following set of results are intended to be a guide to prove result similar to the above in the case of continuously distributed delays An analogue of Lemma I is the following.

LEMMA 3 Let $K_{ij}(s)$,s∈[0,∞), $i,j = 1,2$,denote nonnegative continuous functions integrable on [0,∞) normalised such that

$\int_0^\infty sk_{ij}(s)ds< \infty$; $\int_0^\infty k_{ij}(s)ds=1$, I,j=1,2.

Consider the intergrodifferential system

$$\frac{dY_1(t)}{dt} = -m_{11}Y_1(t) + b_{11}\int_0^\infty k_{ij}(s)\, Y_1(t-s)ds$$

$$-m_{12}\int_0^\infty k_{12}(s)\, Y_2(t-s)ds$$

$$\frac{dY_2(t)}{dt} = -m_{22}Y_2(t) + m_{21}\int_0^\infty k_{21}(s)\, Y_1(t-s)ds$$

$$+b_{22}\int_0^\infty k_{22}(s)\, Y_2(t-s)ds$$

*Where* $m_{ij}, b_{ij}$ $(i,j = 1,2)$ *are positive constants Then*

$$\frac{d|Y_1(t)|}{dt} \leq -m_{11}|\,Y_1(t)\,| + b_{11}\int_0^\infty k_{11}(s)|Y_1(t-s)|ds,$$

$$+m_{12}\int_0^\infty k_{12}(s)|\,Y_2(t-s)|ds,$$

$$\frac{d|Y_2(t)|}{dt} \leq -m_{22}|\,Y_2(t)\,| + m_{21}\int_0^\infty k_{21}(s)|Y_1(t-s)|ds,$$

$$+b_{22}\int_0^\infty k_{22}(s)|\,Y_2(t-s)|ds,$$

*Where* $d|Y_1(t)||/dt$ *is defined by* (A.3).

Proof is again by direct verification.

The proof of the following lemma is similar to that of Lemma 2 above.

Lemma 4. *Consider (A.14) and the following intergrodifferential system:*

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1(t) + b_{11}\int_0^\infty k_{11}(s)Z_1(t-s)ds$$

$$+m_{12}\int_0^\infty k_{12}(s)Z_2(t-s)ds,$$

$$\frac{dZ_1(t)}{dt} = -m_{11}Z_1(t) + b_{11}\int_0^\infty k_{11}(s)Z_1(t-s)ds$$

$$+b_{22}\int_0^\infty k_{22}(s)Z_2(t-s)ds,$$

Together with the initial conditions

$$\infty > Z_i(s) > |Y_i(s)| for\ s \in (-\infty, 0).$$

(A.16)

Then $|Y_i(t)| < Z_i(t) for\ t \geq 0, i = 1,2.$

The following lemma will be useful to prove a result analogous to that of Theorem 3.1 corresponding to a class of continuously distributed time delays.

LEMMA 5. Let $Z_1(t), Z_2(t)$ be the solutions of (A.15)- (A.16) where $(Y_1\ Y_2)$ corresponds to an arbitrary but fixed solutions of (A.13). Then as a consequence of Lemma & it will follow that

$$0 \leq |Y_i(t)| \ Z_i(t) \to \infty \ as\ t \to \infty, \qquad i = 1,2,$$

(A.17)

Whenever

$$\begin{cases} M_{11} > b_{11} + m_{21}, \\ M_{22} > b_{22} + m_{12} \end{cases}$$

PROOF. Define a Lyapunov –like function H where

$$H(Z_1,Z_2)(t) = Z_1(t) + Z_2(t) + b_{11}\int_0^\infty k_{11}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$$

$$+ m_{12}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$$

$$+ m_{21}\int_0^\infty k_{12}(u)\left(\int_{t-u}^t Z_1(s)ds\right)du$$

$$+ b_{22}\int_0^\infty k_{22}(u)\left(\int_{t-u}^t Z_2(s)ds\right)du$$

Where $(Z_1,Z_2)$is any solution of (A.15)-(A.16) for arbitrary $(Y_1,Y_2)$defined on $(-\infty, 0]$ Computing the rate of change of H in (A.19)along the solution of (A.15)we derive that (on using (A.12))

$$\left(\frac{dH(Z_1Z_2(t)}{dt} = [Z_1(t)\{m_{11} - (b_{11} + m_{21})\} + Z_2(t)\{m_{22} - (b_{22} + m_{12})\}\right]$$

$< 0$(using (A.18)) for $Z_1(t) > 0\ and\ Z_2(t) > 0.$

By Lemma 6 below the result will follow

Lemma 6. Let H(t)=H$(Z_1, Z_2)(t)$ be as defined in (A.10) Suppose

$$m_{11} > b_{11} + m_{21} \text{ and } m_{22} > b_{22} |m_{12}$$

Then

$$Z_1(t) + Z_2(t) \to 0 \text{ast} \to \infty.$$

PROOF. Suppose (A.21) dose not hold; then there exists at least one sequence

$\{t_n\} \to \infty$ as n $\to \infty$ such that

$$[m_{11} - (b_{11} + m_{21})]\, Z_1(t_n) + [m_{22} - (b_{22} + m_{12})]Z_2(t_n) > \varepsilon \qquad (A.22)$$

For some sufficiently small positive number $\varepsilon$ and hence

$$\frac{dH(t_n)}{dt} < -\varepsilon, n = 1,2,3, \ldots.$$

By the continuity and differentiability of H with respect to t,

$$\mathrm{H}(t_n - \varepsilon) - \mathrm{H}(t_n) = \frac{dH(t^*)}{dt}(-\varepsilon), \quad t_n - \varepsilon < t^* < t_n$$

If $\varepsilon$ is small we will have

$$= \frac{dH(t^*)}{dt}\varepsilon/2$$

So that

$$\mathrm{H}(t_n) \leq \mathrm{H}(t_n - \varepsilon) - \varepsilon^2/2$$

$$\leq H(t_{n-1}) - \varepsilon^2/2$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2)$$

$$\leq H(t_{n-2}) - 2(\varepsilon^2/2):$$

$$\leq H(t_0) - \mathrm{n}(\varepsilon^2/2): \mathrm{t}_0 = 0.$$

(A.12) and Lemma 5 will imply that $H(t_0) < \infty$ for any $t_0 > 0$. If n is sufficiently large it will follow that $H(t_n) < 0$ since $H(t_0) < \infty$, which contradicts the nonnegativity of H for all $t \geq 0$. Hence $Z_1(t) + Z_2(t) \to 0$ as $t \to \infty$ implying the asymptotic stability of the trivial solution of (A.1)(since $Z_1$ dominates $y_1$)

We note the result of Lemma 5 and 6 will remain valid if (A.18) is replaced by the condition $\mu(A) + \|B\| < 0$

Where A and B are defined as in (6.3).

We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities,

apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems

**Gopalsamy**

# Model For Innocuously Benign Delay In A Periodic System

**Abstract**

Sufficient conditions are obtained for the existence of a unique linearly denomination stable positive periodic solution of a system model of two competitions in a periodic environment with time lags in interspecific interactions. It is shown that if the self- regulating intraspecific interaction effects are strong enough and act without time delays then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a periodic domination. An example in finance will make things clear. Corporate finance for one big plant generates lot of accounts of the workers whose salary shall be disbursed to the debit of cash credit account. Here Cash credit account is a resource as also the one to which debits are raised, a la predator. Some of the accounts from which the credits come are demand bills purchase, usance bills discount, vostro and nostro accounts in respect of foreign bills and Letters of credit. In economics inventories are used to produce output, output is sold to produce income which might be saved in part and recycled in investment. Cosmologically take an example of a Black hole which gormandizes matter and radiates energy. Notwithstanding information paradox, one thing is clear that matter is being created in to radiation. We show how the model could be applied to axiomatic predications, postulation alcovishness, phenomenological methodologies, propositional statement, and other problems in cosmology, philosophy, and even in Green House effects. We study the sine qua non properties of such a system. In fact these models can be applied to such unsolved problems like Yang mills Theory and mass gap, and other problems. As we apply the model, we shall go on dilating the position status and the applicability of the model. We shall not touch the electrical engineering problems which might rise questions like what is charge of a Bank. We may say when a customer enters a Bank; we can say the Bank is charged. These factors are still relativistic and we shall make an in-depth study as we progress. Gopalsamy's research was published in Australian Journal of Mathematical Society vide: B 25349=365 It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 22. Introduction

At the outset we assume that the system satisfies Haimovici is caution. The purpose of this article is to establish a set of sufficient conditions for the existence of a unique linearly asymptotically stable positive periodic solution for a nonautonomous system of the form

$$\frac{dx(t)}{dt} = x(t)f_1\big(t, x(t), y(t-\tau_1)\big),$$

**(1.1)**

$$\frac{dy(t)}{dt} = y(t)f_2\,\big(t, x(t-\tau_2), y(t)\big).$$

In the context of confederation dynamics a model system of the form (1.1) means that the interspecific interactions occur with time delays $\tau_1$ 0, $\tau_2$ 0 in a periodic environment, if $f_1$ *and* $f_2$ are assumed to be periodic in $t$ with a common period $\omega$ 0. Note the interactions are inter specific.

Although model systems in periodic domains have been investigated by a large number of authors (for some references see Gopalsamy (4], [5]), periodically varying environments in time delayed model systems of two or more interacting species have not been investigated. The need for Multitype system models in periodic environments and time delayed interactions are evident from the extensive literature of that type relating to single species. While it is known that periodic ecosystems can support more species than uniform environments (Gopalsamy [4]), time delays might destabilize the system leading to a possible loss of one or more species. Furthermore, periodic environments in population dynamics have some relevance to natural selection and evolution; we will consider this aspect briefly at the end.

The principal finding of the following analysis of (1.1) is that, under suitable hypotheses on $f_1$ *and* $f_2$, the system (1.1) has a linearly asymptotically stable unique positive periodic solution for all nonnegative delays $\tau_1, \tau_2$ (while our numerical experiments suggest that the periodic solution of (1.1) is globally asymptotically stable and is independent of the delay parameters $\tau_1, \tau_2$ it is not known how one can prove such a strong result). Since time delays are usually associated with possible destabilization of otherwise stable systems, we can describe the time delays appearing on (1.1) as “harmless”.

It has been recently shown by the author (Gopalsamy [6]) that if the self –regulating negative feedback in each denomination is strong enough and acts without time delays, then time delays of any length in the interspecific interactions cannot destabilize an otherwise stable system in a temporally uniform domain. The result of this paper is thus an extension to a periodic environment of the "harmless" nature of interspecific interaction delays in a temporally constant environment mentioned above.

## 23. Existence and uniqueness of periodic solutions

We will now formulate the problem of existence of periodic solutions of the periodic delay differential system

$$\frac{dx_1(t)}{dt} = x_1(t) f_1\big(t, x_1(t), x_2(t-\tau_{12})\big),$$

$$t > 0,$$

**(2.1)**

$$\frac{dx_2(t)}{dt} = x_2(t) f_2\big(t, x_1(t-\tau_{21}), x_2(t)\big),$$

(Where $\tau_{12}, \tau_{21}$ are nonnegative constants) in a form suitable for the application of fixed point theory of positive operators. We make the following assumptions regarding $f_1$ and $f_2$ in (2.1).

($H_1$) The partial (Frechet) derivatives of $f_1, f_2$ satisfy the following:

$$\frac{\partial f_i}{\partial x_j} < 0, \qquad i, j = 1, 2 \text{ For all } t > -\tau \; where \; \tau = \max(\tau_{12}, \tau_{21}). \qquad \textbf{(2.2)}$$

($H_2$) $f_1(t, x_1, x_2) \equiv f_1(t + \omega, x_1, x_2)$ for all $t > -\tau, x_1 \; \square 0, x_2 \; \square 0$. **(2.3)**

($H_3$) There exist positive constants $b_{il}$, $b_i^u$, $a_{il}$, $a_i^u$ ($l = 1, 2$) such that

$$b_{il} - a_{il} x_i < f_i(t, x_1, x_2) < b_i^u - a_i^u x_i,$$

$$b_{il}/a_{il} < b_i^u / a_i^u, \qquad \textbf{(2.4)}$$

For $i$= 1, 2 and $(x_1, x_2) \in R_+ \times R_+$. If $\tau = 0$ then one can show that (2.1)-(2.4) will imply that the nondelayed periodic system (2.1) will have an asymptotically stable strictly positive periodic solution (for details see Gopalsamy [5]); this result will become a special case of what we obtain in the following.

We will use the following notation:

$$B_\omega^\tau = \{h: [-\tau, \infty) \to R, h \; is \; continuous, h(t + \omega) = h(t), t \; \square -\tau\},$$

$$B = B_\omega^\tau \times B_\omega^\tau,$$

$$K_1 = \{ h \in B_\omega^\tau | h(t) \geq 0, t \in [-\tau, \infty)\},$$

$$K = K_1 \times K_1,$$

$$||h|| = \left\{ \sup_{t \in [0,\omega],} | h(t)|, h \in B_\omega^\tau \right\}$$

$$||x_1, x_2|| = \max\{||x_1||, ||x_2||, (x_1, x_2) \in B\},$$

$$G_1 = \{(x_1, x_2) \in K, ||(x_1, x_2)|| \leq b_l = \min(b_{1l} / a_{1l}, b_{2l}/a_{2l})\}$$

$$G_2 = \left\{(x_1, x_2) \in K, ||(x_1, x_2)|| \geq b^u = \max\left(\frac{b_1^u}{a_1^u}, \frac{b_2^u}{a_2^u}\right)\right\},$$

$\bar{G} =$ closure of $G_2$ *in* $B$ and

$\frac{\bar{G}_2}{G_1} =$ complement of $G_1 in\ \bar{G}_2$; $G_1 \subset G_2$ by **(2.4)**

The cone $K = K_1 \times K_1$ induces a partial order in $B$ and we will write $(x_1, x_2) \leq (y_1, y_2)$ if and only if $x_1 - y_1 \in K_1$. we define an operator $N$ on $B$ as follows

$$N(x_1, x_2)(t) = \begin{matrix} N_1(t, x_1, x_2) \\ N_2(t, x_1, x_2) \end{matrix} \quad \text{for t> } 0, \qquad \textbf{(2.5)}$$

Where

$$N_1(t, x_1, x_2) = \int_0^\omega G_{11}(m, \omega, t-s) g_1\left(m, s, x_1(x) x_2(s - \tau_{12})\right) ds$$

$$+ \int_0^1 G_{12}(m, 0, t-s) g_1\left(m, s, x_1(s), x_2(s - \tau_{12})\right) ds,$$

$$N_2(t, x_1, x_2) = \int_0^\omega G_{22}(m, \omega, t-s) g_2\left(m, s, x_1(s - \tau_{21}), x_2(s)\right) ds \qquad \textbf{(2.6)}$$

$$+ \int_0^1 G_{21}(m, 0, t-s) g_2\left(m, s, x_1(s - \tau_{21}), x_2(s)\right) ds,$$

For $t > 0$, with

$$g_1\left(m, s, x_1(s), x_2(s - \tau_{12})\right) = m x_1(s) + x_1(s) f_1\left(s, x_1(s), x_2(s - \tau_{12})\right),$$

**(2.7)**

$$g_2\big(m,s,x_1(s-\tau_{21}),x_2(s)\big) = mx_2(s) + x_2(s) f_2\big(s,x_1(s-\tau_{21})x_2,(s)\big),$$

$$G_{ll}(m,\omega,t-s) = e^{-m\omega}(1-e^{-m\omega})^{-1}e^{-m(t-s)}, \qquad i=1,2,$$

**(2.8)**

$$G_{ij}(m,0,t-s) = e^{-m(t-s)}, i \neq j, i,j = 1,2,$$

$$N(x_1,x_2)(t) = \begin{cases} N_1(n\omega + t, x_1, x_2) \\ N_2(n\omega + t, x_1, x_2) \end{cases} for\ t \in [-\tau, 0], \qquad \textbf{(2.9)}$$

*m* being a suitably selected positive constant (more will be said below) while *n* is a positive constant such that $n\omega > \tau$.

We note immediately the following from the definition of $N$: *if* $(x_1,x_2) \in \bar{G}/G_1$ then $N(x_1,x_2)(t)$ ∈ and if $N(x_1,x_2)$ is periodic $\omega$ for $t$ ∈ then $N(x_1,x_2)$ is periodic in *t* with period $\omega$ for $t$ ∈ $-\tau$. The existence of at least one positive periodic solution of (2.1)-(2.4) is established by means of a few lemmas and a result of Schmitt [9].

Lemma 1. *A necessary and sufficient condition for* $(x_1(t),x_2(t))$ , $t$ ∈ $-\tau$ *to be a periodic solution of (2.1)-(2.4) of period* $\omega$ *is that* $(x_1,x_2)$ *be a fixed point of N such that*

$$N(x_1,x_2)(t) = \big(x_1(t),x_2(t)\big) \quad and\ (x_1,x_2) \in K \cap \left(\frac{\bar{G}}{G_1}\right).$$

PROOF. The details of proof are similar to those of Lemma 2 of Busenberg and Cooke [1]. We first verify the continuity of $N(x_1,x_2)$ assuming that $N(x_1,x_2)(t)$ is periodic in *t* of period $\omega$. The continuity of $N(x_1,x_2)$ for $t > 0$ follows from the definition in (2.5)-(2.8) WHICH ALSO PROVIDES $\lim_{t\to 0+} N(x_1,x_2)(t)$ . It will follow from (2.9) and the assumed periodicity of $N(x_1,x_2)$ *that* $N(x_1,x_2)$ is continuous

For $t$ ∈ $-\tau$. Now we will verify the periodicity of $N(x_1,x_2)$ in $t$. A simple computation leads to

$$N_1(t+w,x_1,x_2) - N_1(t,x_1,x_2) = e^{-mt}F_1(t),$$

**(2.10)**

Where

$$F_1(t) = \int_0^{\omega} G_{11}\,(m, \omega, t-s\,)g_1\big(m, s, x_1(s), x_2(s-\tau_{12})\big)ds$$

$$+\int_0^{\omega} G_{12}(m, 0, t-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$+\int_{\omega}^{\omega+t} G_{12}(m, 0, \omega-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$-\int_0^{\omega} G_{11}(m, \omega-s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

$$+\int_0^{t} G_{12}(m, 0, -s)g_1\,\big(m, s, x_1(s), x_2(s-\,\tau_{12})\big)ds$$

**(2.11)**

Which on simplification using the periodicity of $f_1$ shows that $F_1(t) \equiv 0$ whenever $x_1, x_2, \in B$, thus implying the periodicity of $N_1$ in $t$ of period$\omega$. A similar argument for $N_2$ proves that $(x_1, x_2) \in B \Rightarrow N\,(x_1, x_2) \in B$.

Now suppose $(x_1, x_2)$ is a fixed point of *N and let* $(x_1, x_2)\ \in K \cap (\bar{G}/G_1\,)$. Then it is easily seen that $(x_1, x_2)$ is a solution of (2.1) and since $(x_1, x_2)\ \in K \cap (\bar{G}/G_1\,)$ it will follow that $(x_1, x_2)$ is a periodic solution of (2.1).

Conversely, let $(x_1, x_2)$ be a periodic solution of period $\omega$ of (2.1). Then $(x_1, x_2)$ is necessarily a periodic solution of

$$\frac{dx_1(t)}{dt} + mx_1(t) = g_1\,\big(m, t, x_1\,(t), x_2(t-\tau_{12})\big)$$

**(2.12)**

$$\frac{dx_2(t)}{dt} + mx_2(t) = g_2\,\big(m, t, x_1\,(t-\tau_{21}), x_2(t)\big),$$

For any arbitrary real constant *m* and hence for any fixed constant $m > 0$. It follows from (2.12) that any periodic solution of (2.1) and hence of (2.12) is a fixed point of $N(x_1, x_2)$. We claim that such a fixed point is in $K \cap (\bar{G}_2/G_1)$. Suppose $(x_1, x_2) \notin K \cap (\bar{G}_2/G_1)$. the following possibilities can arise:

(xxii) Either $x_1, x_2$ or both $x_1$ *and* $x_2$ may become negative for some value of *t* and it will follow from (2.1) that subsequent value of $x_1, x_2$ or $x_1$ *and* $x_2$ will be negative and hence from (2.4) and (2.1) their derivatives will be negative for all subsequent *t* , implying that $x_1, x_2$ or $x_1$ *and* $x_2$ cannot be periodic.

(xxiii) While $(x_1, x_2) \in K \subset B$ there may be a value of *t*, say $t^*$, such that $x_1(t^*) > b_1^u/a_1^u$ or $x_2(t^*) > b_1^u/a_1^u$ and in such a case $x_1(t) or\ x_2(t)$ cannot be periodic since respectively $dx_1(t^*)/dt < 0\ or\ dx_2(t^*)/dt < 0,$ which will imply that the values of $x_1(t^*)\ or\ x_2(t^*)$ cannot be repeated from below.

(xxiv) Similarly if there is a $t^*$ such that $0 < x_1(t^*) < b_{1l}/a_{1l}\ or\ 0 < x_2(t^*) < b_{2l}/a_{2l}$ then such solutions cannot be periodic since $dx_1(t^*)dt > 0$ or $dx_2(t^*)dt > 0$, which will mean that the values of $x_1(t^*)$ or $x_2(t^*)$ cannot be repeted from above.

Thus every periodic solution of (2.1)-(2.4) is a fixed point of *N* in $K \cap \left(\frac{G_2}{G_1}\right) \subset B$.

**LEMMA 2.** *Let N be as in (2.5)-(2.6). Suppose* $\lambda \in (0,1) and\ \lambda N(x_1, x_2) = (x_1, x_2)$ *for some* $((x_1, x_2) \in K - \{0\}$. *then there exist* $\alpha_1^* > 0, \alpha_2^* > 0$ *such that*

$$0 < x_1(t) \leq \alpha_1^* < b_1^u / a_1^u,$$

$$0 < x_2(t) \leq \alpha_2^* < b_2^u / a_2^u,$$

*for* $t \in [-\tau, \infty)$.

PROOF. If $\lambda N(x_1, x_2) = (x_1, x_2)$ then from the definition of N we get for all $t > 0$ and $(x_1, x_2) \in K - \{0\}$,

$$\frac{dx_1(t)}{dt} = \lambda\, x_1\,(t) f_2(t, x_1(t), x_2\,(t-\tau_{12})) + x_1\,(t) m(\lambda - 1)$$

**(2.13)**

$$\frac{dx_2(t)}{dt} = \lambda\, x_2\,(t) f_2(t, x_1(t-\tau_{21}), x_2\,(t)) + x_2\,(t) m(\lambda - 1)$$

For all $t$ such that $x_1(t) = b_1^u/a_1^u$ or $x_2(t) = b_2^u/a_2^u$ we have $\frac{dx_1}{dt} < 0$ or $\frac{dx_2}{dt} < 0$; but these are impossible since $x_1, x_2$ are periodic in t. Thus by the continuity and periodicity of $x_1, x_2$ we have $x_1(t) \leq x_1^* < b_1^u/a_1^u$, $x_2(t) \leq x_2^* < b_2^u/a_2^u$ for all $t \geq 0$; by the definition and periodicity the same will follow for all $t \geq -\tau$. A similar argument holds when $x_1(t) < 0$ or $x_2(t) < 0$, noting that in such a case $dx_1/dt > 0$ or $dx_2/dt > 0$, again contradicting the periodicity of $(x_1, x_2)$ in *t*.

LEMMA 3. *The operator N maps* $K \cap (G_2/G_1)$ *in to K and is completely continuous*

PROOF. From the definition of N and (2.4) we derive that

$$N_1(t, x_1, x_2) = \int_0^\omega G_{11}\,(m, \omega, t-s) x_1(s)[m + b_{1l} - a_{1l} x_1(s)] ds$$

$$+ \int_0^\omega G_{12}\,(m, 0, t-s) x_1(s)[m + b_{1l} - a_{1l} x_1(s)] ds.$$

**(2.14)**

For $(x_1, x_2) \in K \cap (G_2/G_1)$, we have $0 < x_1 < b_1^u/a_1^u$ and hence we get $N_1(t, x_1, x_2) > 0$ for all $t \geq 0$ by a proper choice of the positive constant say $m_1$.

We make a similar choice say $m_2$ for *m* for the case $N_1(t, x_1, x_2) \geq 0$ and then choose *m* = max $(m_1, m_2)$. Thus we have

$$N: K \cap (G_2/G_1) \to K \qquad for\ t \geq 0,$$

And by the periodicity of N we have the same for $t - \tau$.

To prove the complete continuity of N we proceed as follows. We know that if $(x_1, x_2) \in B$ then $N(x_1, x_2) \in B$ and so we can restrict $t$ to the set $[0, n\omega], n\omega > r$. We write $N = L \circ Q$ where $L$ is the linear operator, $L : B \rightarrow B$, $L(u_1, u_2)(t)$

$$= \int_0^{\omega} G_{11}(m, \omega, t-s) \begin{bmatrix} u_1 & (s) \\ u_2 & (s) \end{bmatrix} ds, + \int_0^{t} G_{12}(m, 0, t-s) \begin{bmatrix} u_1 & (s) \\ u_2 & (s) \end{bmatrix} ds$$

**(2.15)**

And Q is the nonlinear operator

$$Q(u_1, u_2)(t) = \begin{matrix} g_1(m, t, u_1(t), u(t-\tau_{12})) \\ g_2(m, t, u_1(t-\tau_{21}), u_2(t)) \end{matrix} for \ t > 0,$$

**(2.16)**

$$Q(u_1, u_2)(t) = Q\,(u_1, u_2)(n\omega + t) \quad for\ t \in [-\tau, 0].$$

It can be found that

$$||L(u_1, u_2)(t)|| \leq max \left\{ \frac{m+b_1^u}{m}, \frac{m+b_2^u}{m} \right\} ||u_1, u_2||,$$

**(2.17)**

$$\left\| \frac{d}{dt} L(u_1, u_2) \right\| \leq 2\ max \left\{ \frac{m+b_1^u}{m}, \frac{m+b_2^u}{m} \right\} ||u_1, u_2||.$$

The complete continuity of the linear map $L$ follows from (2.17). Now since Q is continuous and maps bounded sets into bounded sets it will follow that $L \circ Q$ is completely continuous.

To prove the existence of periodic solution of (2.1)-(2.4) we will use the following result of Schmitt [9].

**THEOREM (Schmitt).** *Let $H_1$ and $H_2$ be nonempty bounded open neighborhoods of the origin in a Banach space B with a positive cone K such that $H_1 \subset H_2$ and let N: $K \cap (H_2 / H_1) \to K$ be completely continuous and which satisfies the following conditions:*

***(o)*** *There exists $(z_1, z_2) \in K \subset B, \|(z_1, z_2)\| = 1$ and a $\mu$ satisfying*

$$\mu > sup\{\|(x_1, x_2) - N(y_1, y_2)\|, (x_1, x_2) \in K \cap H_1, (y_1, y_2) \in K \cap \partial H_2\}$$ *such that all solution* $(x_1, x_2) \in K \cap (H_2/H_1)$ *of* $(x_1, x_2) = N(x_1, x_2) + \lambda(z_1, z_2) with\ 0 < \lambda < \mu\ satisfy\ (x_1, x_2) \notin \partial H_1.$

*(p) All solution* $(x_1, x_2) \in K \cap (H_2/H_1)\ of\ (x_1, x_2) = \lambda N(x_1, x_2), 0 < \lambda < 1\ satisfy\ (x_1, x_2) \notin \partial H_2.$

*Then N has at least one fixed point in* $K \cap (H_2/H_1)$.

We can now present our existence result.

**THEOREM 1.** *Assume that* $f_1, f_2$, *of* (2.1) *satisfy* (2.2)-(2.4). *then there exists at least one positive periodic solution of* (2.1)-(2.4).

PROOF. By Lemma 1, it is enough to show that the operator $N: K \cap (G_2/G_1) \to K \subset B$ has at least one fixed point in $G_2/G_1$ and we will do this by applying the result of Schmitt [9]. First we let

$$H_2 = \{(x_1, x_2) \in B | \ ||x_1, x_2|| < \max(b_1^u/a_1^u, b_2^u/a_2^u)\} \equiv G_2,$$

$$H_1 = \{(x_1, x_2) \in B | \ ||x_1, x_2|| < \min(b_{1l}/a_{1l}, b_{1l}/a_{1l})\} \equiv G_1.$$

By Lemma 3, $N: K \cap (H_2/H_1) \to K$ is completely continuous. Choose $\mu = \infty$ and $(z_1, z_2) \equiv (1.1)$. Then all solution $(x_1, x_2) \in K \cap (H_2/H_1)$ of $(x_1, x_2) = N(x_1, x_2) + \lambda(z_1, z_2), 0 < \lambda < \infty$, satisfy (2.1) and hence by their periodicity $(x_1, x_2) \notin \partial H_1$. The condition (a) of Schmitt's theorem is satisfied since $\mu = \infty$ by choice and $(x_1, x_2)$ and $(z_1, z_2)$ are periodic. By Lemma 2 the condition (b) of Schmitt's theorem is satisfied. Thus there exists at least

one fixed point of N in $K \cap (H_2/H_1)$ and hence there exists at least one positive periodic solution of (2.1)-(2.4) and the proof is complete.

We remark that so far no restriction on the magnitudes of the delays $\tau_{12}, \tau_{21}$ has been imposed except that they be nonnegative.

## 24. Uniqueness of the periodic solution

Applications of fixed point theorems to prove the existence of solutions of nonlinear problems usually suffer from the drawback that the uniqueness of solutions has to be proved often by other methods. We will use the following result of Smith [11] which is a slight variation of his Theorem 3.5 (see also Busenberg and Cooke [1]).

THEOREM (Smith). *Let N be an operator on a cone K in the Banach space B and let* $K^* = \{(x_1, x_2) \in K \cap B, 0 \leq x1,x2 \leq C*$ *where* $C*$*some positive constant is. Assume the following:*

(v) If $(x_1, x_2) \in K^* - \{(0,0)\}$ and $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

(w) N is monotonic nondecreasing on $K^*$.

(x) If $(x_1, x_2)$ is in the interior of $K^*$ and if $0 < \lambda < 1$ then there exists an $\eta > 0$ depending on $(x_1, x_2)$ and $\lambda$ such that $N(\lambda x_1, \lambda x_2) \geq \lambda(1 + \eta) N(x_1, x_2)$. Then N has at most one fixed point different from $(0,0)$ in $K^* \subset B$.

The cone K= $K_1 \cap K_1$ where $K_1 = \{h \in B_\omega^\tau : h(t) \geq 0, t \in [-\tau, \infty)\}$ is slightly too wide in that solutions of the form $\{x_1(t), 0\}$ and $\{0, x_2(t)\}$ for (2.1) are also possible. We choose a narrower cone $K^*$ in $B$ as follows: define

$$S = \{(x_1, x_2) \in B : b_{1l}/a_{1l} \leq \|x_1\| \leq b_1^u/a_1^u,\ b_{2l}/a_{2l} \leq \|x_2\| \leq b_2^u/a_2^u,\}, \quad \textbf{(3.1)}$$

And let

$$K^* = \{\lambda(x_1, x_2) \in B : (x_1, x_2) \in S, \lambda \geq 0\}. \quad \textbf{(3.2)}$$

Since $\hat{G}_2/G_1 \subset K^*$, it is not difficult to see that N: $\hat{G}_2/G_1 \to B$ has at least one fixed point which is a periodic solution of (2.1)-(2.4). We will prove that this solution is unique by applying the result of Smith cited above. The following lemmas are needed first.

LEMMA 4. If $(x_1, x_2) \in K^* - \{0.0\}$ and if $N(x_1, x_2) = (x_1, x_2)$ then $(x_1, x_2) \in$ interior of $K^*$.

PROFF. If $(x_1, x_2) = N(x_1, x_2)$ and $(x_1, x_2) \in K^* - \{0,0\}$, then $(x_1, x_2)$ is a periodic solution of (2.1) –(2.4) and hence we have

$$(b_{1l} - a_{1l}x_1)x_1 < \frac{dx_1}{dt} < (b_1^u - a_1^u x_1)x_1, \qquad (\mathbf{3.3})$$
$$(b_{2l} - a_{2l}x_2)x_2 < \frac{dx_2}{dt} < (b_2^u - a_2^u x_2)x_2,$$

And hence any nonzero periodic solution $(x_1, x_2)$ of (2.1)-(2.4) has to be such that

$$b_{1l}/a_{1l} < x_1(t) < b_1^u - a_1^u,$$

$$b_{2l}/a_{2l} < x_2(t) < b_2^u - a_2^u,$$

For $t > -\tau$, which proves the lemma.

LEMMA 5. The mapping $N$ defined by (2.5)-(2.9) is monotone nondecreasing on $K^*$.

PROOF. Let $(x_1, x_2) \in K^*$, $(y_1, y_2) \in K^*$ such that $(x_1 - y_1, x_2 - y_2) \in K^*$ and $(x_1, x_2) \neq (y_1, y_2)$. Then $x_1 \geq y_1$, $x_2 \geq y_2$. From the definition of $N$ we derive that

$N_1(t, x_1, x_2) - N_1(t, y_1, y_2)$

$$= \int_0^\omega G_{11}(m, \omega, t-s)\{m[x_1(s)] + x_1(s)f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s)f_1(s, y_1(s), y_2(s-\tau_{12}))\}ds$$

$$+ \int_0^t G_{12}(m, 0, t-s)$$

$$x\{m[x_1(S) - y_1(s)] + [x_1(s)f_1(s, x_1(s), x_2(s-\tau_{12})) - y_1(s)f_1(s, y_1(s), y_2(s-\tau_{12}))]\}ds \qquad \mathbf{(3.4)}$$

With a similar relation for $N_2(t, x_1, x_2) - N_2(t, y_1, y_2)$. Since $f_1$ is continuous, $f$ is bounded on $K^*$; also it will follow from $(x_1, x_2) \neq (y_1, y_2)$ that

$$0 \not\equiv x_1(s) - y_1(s) \geq 0 \text{ and } 0 \not\equiv x_2(s) - y_2(s) \geq 0;$$

A sufficiently large positive constant $m$ will make the right side of (3.4) positive.A similar argument will lead to $N_2(t, x_1, x_2) - N_2(t, y_1, y_2) > 0$ perhaps with a larger constant $m$. Finally we can choose the bigger of these two constants $m$ to yield

$$N(x_1, x_2) - N(y_1, y_2) \text{ for } (x_1 - y_1, x_2 - y_2) \in K^* - \{(0,0)\}.$$

The monotonicity of $N$ on $K^*$ follows.

LEMMA 6. Let $(x_1, x_2) \in$ interior of $K^*$ and let $\lambda$ be $a$ constant such that $0 < \lambda < 1$. Suppose $f_1, f_2$ satisfy (2.2) - (2.4). Then there exists $a$ Positive constant $\eta$ such that

$$N(\lambda x_1, \lambda x_2) \geq \lambda(1+\eta) N(x_1, x_2). \quad \textbf{(3.5)}$$

PROOF. A direct computation leads to

$$N_1(t, \lambda x_1, \lambda x_2) - \lambda N_1(t, x_1, x_2)$$

$$= \lambda \int_0^{\omega} G_{11}(m, \omega, t-s) g_1\big(s, x_1(s), x_2(s-\tau_{12})\big) H_1(s, \lambda, x_1, x_2) ds$$

$$+ \lambda \int_0^{t} G_{11}(m, 0, t-s) g_1\big(s, x_1(s), x_2(s-\tau_{12})\big) H_1(s, \lambda, x_1, x_2) ds$$

**(3.6)**

Where

$$H_1(s, \lambda, x_1, x_2) = \frac{m\, x_1(s) + x_1(s) f_1\big(s, \lambda\, x_1(s), \lambda\, x_2(s-\tau_{12})\big)}{m\, x_1(s) + x_1(s) f_1\big(s, x_1(s), x_2(s-\tau_{12})\big)} - 1. \quad (\mathbf{3.7})$$

Since $f_1$ is decreasing (see (2.2)) with respect to $x_1$ and $x_2$ there exists a positive constant say $\eta_1$ such that

$$H_1(s, \lambda, x_1, x_2) \geq \eta_1 \text{ for } 0 < \lambda < 1, s > 0, (x_1, x_2) \in K^*, \quad \textbf{(3.8)}$$

And hence

$N_1(t, \lambda x_1, \lambda x_2) \geq \lambda(1+\eta_1)N_1(t, x_1, x_2)$.

A similar argument will show that there exists an $\eta_2 > 0$ such that

$N_2(t, \lambda x_1, \lambda x_2) \geq \lambda(1+\eta_2)N_2(t, x_1, x_2)$,

And the result follows if we choose $\eta = \min(\eta_1, \eta_2)$.

Now the following uniqueness result is a direct consequence of the result of Smith [11] and the Lemmas 4, 5, 6.

THEOREM 2. *Suppose $f_1, f_2$ of (2.1) satisfy (2.2) – (2.4). Then there exists a unique $\omega$ - periodic strictly positive periodic solution of (2.1).*

## 4. Linear stability of the periodic solution

Let $\{x_1^*(t), x_2^*(t), t \geq -\tau\}$ be the strictly positive unique periodic solution of (2.1). Let $\{x_1(t), x_2(t), t \geq -\tau\}$ be any other solution of (2.1). To study the linear stability of $(x_1^*, x_2^*)$ we let

$$\begin{aligned} x_1(t) &\equiv x_1^*(t)\{1 + X_1(t)\}, \\ x_2(t) &\equiv x_2^*(t)\{1 + X_2(t)\}, \end{aligned} \quad t \geq -\tau, \tag{4.1}$$

And derive that the perturbations $(x_1, x_2)$ are governed by the linear variational system (after neglecting the non linear terms in $x_1, x_2$)

$$\begin{aligned} \frac{dX_1(t)}{dt} &= -a_{11}(t)\, x_1^*(t)X_1(t) - a_{12}(t)x_2^*(t-\tau_{12})\, X_2(t-\tau_{12}) \\ \frac{dX_2(t)}{dt} &= -a_{21}(t)x_1^*(t-\tau_{21})x_1(t-\tau_{21}) - a_{22}(t)x_2^*(t)x_2(t), \end{aligned} \tag{4.2}$$

For $t > 0$, where

$$\left.\begin{aligned} -a_{11}(t) &= \frac{\partial f_1}{\partial x_1} \\ -a_{12}(t) &= \frac{\partial f_1}{\partial x_2} \end{aligned}\right\} \text{Evaluated at } (t, x_1^*(t), x_2^*(t-\tau_{12})), \tag{4.3}$$

$$\left.\begin{aligned} -a_{22}(t) &= \frac{\partial f_2}{\partial x_1} \\ -a_{22}(t) &= \frac{\partial f_2}{\partial x_2} \end{aligned}\right\}$$ Evaluated at $(t, x_1^*(t-\tau_{21}), x_2^*(t))$.

To establish the linear asymptotic stability of $(x_1^*(t), x_2^*(t))$ it is enough to prove that of the trivial solution of (4.2) - (4.3) We will do this by means of a combination of a comparison differential system and a Lyapunov function. First we define

$$\frac{d|X_1(t)|}{dt} = \begin{cases} \frac{dX_1(t)}{dt} & if\ X_i(t) \geq 0, \\ -\frac{d\setminus X_1(t)}{dt} & if\ X_i(t) < 0, \end{cases} \quad i = 1,2, \qquad \textbf{(4.4)}$$

And derive from (4.2)- (4.3) the following comparison system of differential inequalities:

$$\begin{aligned} \frac{d|X_1(t)|}{dt} &\leq -b_{11}, x_1^*(t)\,|X_1(t)| + b_{12}x_2^*(t-\tau_{12})\,|X_2(t-\tau_{12})|, \\ \frac{d|X_2(t)|}{dt} &\leq -b_{21}, x_1^*(t-\tau_{21})\,|X_1(t-\tau_{12})| - b_{22}x_2^*(t)\,|X_2(t)|, \end{aligned} \qquad \textbf{(4.5)}$$

For $> 0$, where

$$-b_{11} = \max_{t\in-\tau}\{-a_{11}(t)\},$$

$$\begin{aligned} b_{12} &= |\min_{t\in-\tau}\{-a_{12}(t)\}| \\ b_{21} &= |\min_{t\in-\tau}\{-a_{21}(t)\}| \end{aligned} \qquad \textbf{(4.6)}$$

$$-b_{22} = \max_{t\in-\tau}\{-a_{22}(t)\}$$

We can now formulate the result as follows. By the periodicity of $a_{ij}$, the constants $b_{ij}$ are independent of $\tau$ (see also (2.2)).

THEOREM 3. *The unique strictly positive periodic solution* $\{x_1^*(t), x_2^*(t), t \in -\tau\}$ of $(2.1) - (2.4)$ *is linearly asymptotically stable provided*

$$|-b_{11}| > b_{21} \text{ and } |-b_{22}| > b_{12},$$

**(4.7)**

*Where the constants* $b_{ij}$ $(\mathrm{i,j} = 1,2)$ *are as defined in* (4.6).

PROOF. Consider the scalar valued Lyapunov-like function $v(t) = V(t, X_{1,}X_{2,})$where

$$v(t) = V(t, X_{1,}X_{2,}) = |X_{1,}(t)| + |X_{2,}(t)| \int_{t-\tau_{12}}^{t} b_{12} x_2^*(u)|X_{2,}(u)|du$$

$$+\int_{t-\tau_{21}}^{t} b_{21} x_1^*(u)|X_{1,}(u)|\mathrm{du}. \qquad \textbf{(4.8)}$$

Since $x_1^*, x_2^*$ are strictly positive, it is found from (4.7)-(4.) that

$$\frac{dv}{dt} \leq -[(|b_{11}| - b_{21})\, x_1^*(t)|X_{1,}(t)| + (|b_{22}| - b_{22})\, x_2^*(t)|X_2(t)|] \qquad \textbf{(4.9)}$$

For $t > 0$. It follows from (4.7)-(4.9) that $V(t, X_{1,}X_{2,}) \geq |X_{1,}(t)| + X_{2,}(t)|$ and $dv/dt \leq 0$,which will imply the stability of the trivial solution of (4.2)-(4.3). We claim that (4.9) also implies the asymptotic stability of the trivial solution of (4.2)-(4.3)in the sense that

$$\lim_{t-\infty} [\,|X_{1,}(t)| + |X_{2,}(t)|\,] = 0,$$

**(4.10)**

For all bounded initial conditions $X_{1,}(s), X_{2,}(s), s \in \{\mathcal{T}, 0\}$. $Suppose (4.10) is\ not$ true Then we can find an increasing sequence $\{t_m\} \to \infty$ as $m \to \infty$ such that

$$(|b_{11}| - b_{21})\, x_1^*(t_m)|X_{1,}(t_m)| + (|b_{22}| - b_{12} x_2^*(t_m)|X_2(t_m)| > \varepsilon,$$

M=1, 2, 3…

**(4.11)**

Using the right derivative *dv/dt* of $v$ we derive that

$$v(t_m - \varepsilon) - \mathrm{v}(t_m) \leq -\int_{t_m-\varepsilon}^{t_m} \frac{\mathrm{dv}(\mathrm{T})}{\mathrm{dT}}\,\mathrm{d}\tau.$$

**(4.12)**

If $\varepsilon$ is small enough we have from (4.9)(4.11),(4.12)that $dv(\tau)/dr < (-\varepsilon/2)$ and

$$v(t_m - \varepsilon) - \mathrm{v}(t_m) \geq \varepsilon^2/2 \text{ and hence}$$

$$v(t_m) \geq v(t_m - \varepsilon) - (\varepsilon^2/2)$$

$$\geq v(t_m - 1)\text{-}(\varepsilon^2/2)$$

$$\geq v(t_m - 2)\text{-}(\varepsilon^2/2)$$

$$\geq v(t_0)\text{-m}\,(\varepsilon^2/2)$$

Thus for sufficiently large $m, v(t_m)$becomes negative which contradicts the Nonnegative nature of $v\ for\ t \geq 0$.

this contradiction shows that $|X_{1,}(\mathrm{t})| + |X_2(\mathrm{t})| \to 0$ as $\mathrm{t} \to \infty$ and the proof is complete.

As an example let us consider the case of two species Lotka- Volterra type competition in a periodic environment for for which we have

$$f_1\big(\mathrm{t}, x_1(\mathrm{t}), x_2(\mathrm{t} - T_{12})\big)1 \quad \left\{ \frac{x_1(\mathrm{t}) + a x_2(\mathrm{t} - T_{12})}{K_1(\mathrm{t})} \right\}$$

$$f_2(\mathrm{t}, x_1(\mathrm{t} - {}_{21}) x_2(\mathrm{t})) = 1 - \left\{ \frac{\beta x_1(\mathrm{t} - T_{21}) + x_2(\mathrm{t})}{K_2(\mathrm{t})} \right\}$$

Where $\alpha, \beta$ are positive constants and $k_{1,} k_{2,}$ are strictly positive continuous periodic functions with a common period Equation (4.6)and (4.7) for this special case lead to

$$\left|\operatorname{axm}_{t>T}\{-a_{11}(t)\}\right| = \max_{t>T}(1/K_1(t)) > \min_{t>T}\left(-a_{21}(\mathrm{t})\right)\Big| = \min_{t\in}, \beta/K_2(\mathrm{t})$$

$$\max_{t\in T}\{-a_{22}(t)\}\Big| = \max_{t\in T}(1/K_2(t)) > \min_{t\in T}\left(-a_{12}(\mathrm{t})\right)\Big| = \min_{t\in} \mathfrak{a}/K_1(\mathrm{t})$$

The inequalities reduce in the absence of temporal variations in $k_{1,}k_{2,}$ to the inequality 0< $\mathfrak{a}$,$\beta$<1, which is well known for the global asymptotic stability of the positive steady state corresponding to the above example in a constant systemic domain the above example shows also that delays in interspecific interactions have no influence on the local asymptotic stability of the unique periodic solution of the periodic ecosystem A number of cases of this type in a constant environment have been established recently by the author (Gopalsamy [6]).A numerical example is illustrated at the end.

**Consequences**

It can be found that all the conclusions on the existence uniqueness and stability of a periodic solution of (2.1)-(2.4) are valid when $OT_{12}$=0 or $T_{12}$=0 or both $T_{12}$=0 and $T_{21}$=0: thus it will follow that delays in interspecific competition do not destabilize an otherwise stable coexistence (oscillatory or nonoscillatory) of two competing types of interacting particles atoms molecules transactions interaction. In forthcoming work we will extend the analysis of the system (2.1)-(2.4) to a system of n-competing species with a hypothesis similar to that of (2.4) It is however an open problem to investigate the effects of delays in intraspecific competition in a periodic milieu

Pianka ([8],pages 139-140) remarks that “In real population rate of increase competitive abilities and carrying capacities do very from individual to individual with population density in space and time Indeed temporal variation in the environment may often allow coexistence by continually altering the competitive abilities of population inhabiting it Moreover time lags are doubt lessly of some importance in real population Finally a heterogeneous environment may allow real competitors to evolve divergent resource utilization patterns and to reduce interspecific competitive inhibition

To understand the mechanism of selection pressure against competitive inhibition we first write the Malthusian parameters $\mu_1$, $\mu_2$ of fitness for the two species as follows (see MacArthur[7],page 226-230):

$$\mu_1=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_1\left(t,x_1(t),x_2(t-T_{12})\right)$$

$$\mu_2=\frac{1}{x_1(t)}\frac{dx_1(t)}{dt}=f_2(t,x_1(t-T_{21}),x_2(t))$$

Since $\left(x_1(t),x_2(t)\right)\rightarrow\left(x_1^*(t),x_2^*(t)\right)$ for large t we can write the parameters of finess in the from

$$\mu_1=f_1\left(t,x_1^*(t),x_2^*(t-T_{12})\right)$$

$$\mu_2=f_2\left(t,x_1^*(t-T_{12}),x_2^*(t)\right)$$

In these relations we can use

$$x_2^*(t-T_{12}),x_2^*(t)-T_{12}\frac{dx_2^*(t-T_{12}+\theta_1T_{12}}{dt}$$

$$x_1^*(t-T_{12}),x_1^*(t)-T_{21}\frac{dx_1^*(t-T_{21}+\theta_2T_{21}}{dt}$$

$\theta_1,\theta_2,\in(0,1)$ and rewrite $\mu_1,\mu_2$ in the form

$$\mu_1=f_1\left(t,x_1^*(T),x_2^*(t)-T_{12}\frac{dx_2^*(t-T_{12}+\theta_1T_{12}}{dt}\right.$$

$$\mu_2=f_2\left(t,x_1^*(T),-T_{21}\frac{dx_1^*(t-T_{21}+\theta_2T_{21}}{dt},x_2^*(t)\right)$$

$\theta_1,\theta_2,\in(0,1)$

It is seen that the presence of the delays $T_{12}$ and $T_{21}$ provides additional sources of variability in the finess parameters and thereby in the selection pressure against interspecific competition Since a necessary and sufficient condition for natural selection to act is differential (variable) rates of confederation growth we find that time delays coupled with periodicity of an domain may allow real to evolve divergent resource utilization patterns and to reduce

Interspecific competitive inhibition for some remarks relating to the significance of a periodic environment to sustain the coexistence of a large number of confederation types in the absence of delays we refer to a recent article (Gopalsamy)The following example has been numerically solved and the solutions are displayed graphically for two different values of the delay parameter Ƭ:

$\frac{dx_1(t)}{dt} x_1(t)[1 - x_1(t) + 0.4x_2(t - \mathcal{T})/k_1(t)]$

$\frac{dx_2(t)}{dt} x_2(t)[1 - \{0.8\, x_1(t - \mathcal{T}) + x_2(t)\}/k_2(t)]$

Where $k_1(t) = 20 + \sin(\pi t)$ and $k_2(t)$=2.5+sin ($\pi$t)

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There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies

systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational

content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification.

## GOPALSAMY'S MODEL APPIED TO COSMOLOGICAL PROBLEMS

## GLOBAL ASYMPTOTIC STABILITY IN AN ALMOST PERIODIC LOTKA- VOLTERRA SYSTEM

**Abstract**

**Gopalsamy studied** Sufficient conditions are obtained for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of a Lotka- Volterra system with almost periodic coefficients. Research was published Australian. Math Soc. B.27 (1986) 346-360. Gopalaswamy in his Delayed responses and stability in two species systems (Journal of Australian mathematical SocietyB25 473-500) has shown that if intraspecific self- regulating negative feedback effects are strong enough such that a nontrivial steady state of a two species system is locally asymptotically stable, then time delays in the positive feedback as well as in other interspecific interactions cannot destabilize the system and hence delay induced instability leading to persistent oscillations is impossible whatever the magnitude of the time delays. A method is also proposed for an estimate of

decay rate of perturbations. Sole intention, main objective and primary concern of the study is to show how this model has wide ranging amplitudinal ramifications such as in finance, econometrics, cosmology, propositions, phenomenology, philosophy and other fields. In finance cash credit account is an account that is given by a bank to the customer. This forms the resource which is assiduously and widely competed by such factorial representationalities such as bills received in the name of the company, foreign bills received by the company. On the other hand the bills purchased and discounted also form a creditional value for the cash credit account. Thus there are various debits and credits to the Cash credit account. Thus Cash credit account forms a sort of both predator and prey in consideration to the fact they receive both credits and debits. In econometrics, investment leads total output thus depleting part of the investment, not wholly or in full measure and output leads to savings, consumption, and part of the income so obtained could be ploughed back to the investment portfolio. In this connection reference is invited to the authors DLitt thesis on the Policy studies wherein details studies of the application of the predator and prey models with nearly 57 parameters and concomitant 167 equations have been suggested corresponding to an optimization problem and duality of motion. It is the fundamental theorem that energy is always conserved. And energy is matter. In many of the cosmological interactions this holds good. In many of propositional and problem approaches as in the case of philosophers' like Derrida, Kant, Deleuze, and others. One extant example is the case of Horava Lifshitz gravity in which interactions are of the type to which the model could be applied. All the result hold good for confederation of atoms, molecules, creation rate, destruction rate, financial account, cosmological interactions, where energy is transferred from one to another, interaction in gravitational field. Quintessentially population is the field and the creation destruction rates or debit credit rates are counted in the same vein as that of predator prey. Note that the model is applicable to each and every module. Quintessentially each parameter should have been given different name such as superscript 1,2, etc., as that has been done in the foregoing and a formal concatenation done at the end. Due to spatial restrictions we have not been able to do the line of thought and manner of analysis as has been done hereinbefore. This is the most important aspect of the model. It is to be stated in unmistakable terms and unequivocal enucleation that the nine module concatenated model becomes applicable for each and every sentence of the model that has been incorporated, including the equation. For example, if there is an equation A+B=C , then A can be differentiated and written in the governing equation form and the corresponding model equations as is stated in the foregoing modular structure gets changed which could simply be represented by a different subscript and superscript for expression felicity. It is in this manner each term has to be solved and finally the present incorporated structure has to be reproduced for ultimate consolidation and corporation by concatenation. This would take lot space to the extent of thousands of pages and long time for expatiation and exposition itself. It is of seminal, cardinal and paramount importance that this has to be borne in mind before the concatenated models are read with each and every module, repeat each and every module. In fact nine module 54 storey structures had to be inculpated in a concatenated form for each and every paper. This again has been refrained from for the spatial restrictions. Once this fundamental truth is stated and understood explicit, the further reading becomes a cake walk.

## 15. Introduction

At the outset we assume that the system. This article is concerned with the derivation of a set of "easily verifiable" sufficient conditions for the existence of a globally asymptotically stable strictly positive (component wise) almost- periodic solution of the Lotka- Volterra system

$$\frac{dx_1(t)}{dt} = x_i(t)\left\{b_1(t) - \sum_{j=1}^{n} a_{ij}\,(t)x_j(t)\right\};$$

$$t>t_0;\ t_0\ \in (-\infty,\infty);\ i = 1,2,\dots,n, \qquad (1.1)$$

where $b_1, a_{ij}\ (i,j = 1,2,\dots.,n)$ are nonnegative almost periodic functions defined for $t \in (-\infty,\infty)$.

In question satisfies Haimovici's equation based on age wise stratification given in the foregoing. Note that Haimovici's equations can be applied to Lotka-Volterra system term by term.

Let R and $R_n$ respectively denote the set of all real numbers and the $n$ - dimensional real Euclidean space; $R_n^+$ will denote the nonnegative cone of $R_n$. Since almost- periodic functions defined on R are bounded, we define the constants $b_i^l, b_l^u, a_{ij}^l, a_{ij}^u, (i,j = 1,2,\dots,n)$ by the following:

$$\inf_{t\in R} b_l(t) = b_i^l; \qquad \sup_{t\in R} b_l(t) = b_l^u, \qquad i,j = 1,2,\dots,n.$$

$$\inf_{t\in R} a_{lj}(t) = a_{lj}^l; \qquad \sup_{t\in R} a_{lj}(t) = a_{ij}^u.$$

We shall study the almost- periodic system (1.1) with the following assumption on the coefficients of (1.1);

$$b_l^l > 0; \qquad a_{ll}^l > 0, \tag{1.2}$$

$$b_l^l > \sum_{\substack{j=1 \\ j\neq i}}^{n} a_{lJ}^u (b_j^u / a_{iJ}^l). \tag{1.3}$$

Since the solutions of (1.1) corresponding to nonnegative initial conditions remain nonnegative subsequently, it will follow that

$$\frac{dx_1}{dt} \leq x_l\{b_l^u - a_{ll}^l x_i\}; \qquad t > t_0, i = 1,2,\dots,n, \tag{1.4}$$

As a consequence of which we will have

$$0 < x_l(t_0) \leq b_l^u / a_{ll}^l = x_i^u \Rightarrow x_i(t) \leq b_l^u / a_{Il}^l, \qquad i = 1,2,\dots,n; t \geq t_0. \tag{1.5}$$

(1.8) And (1.5) together lead to

$$\frac{dx_1}{dt} \leq x_l \left\{ b_l^u \sum_{\substack{j=1 \\ j\neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right) - a_{ll}^l x_i \right\}, \qquad i = 1,2,\dots,n; \tag{1.6}$$

Now (1.2), (1.3) and (1.6) lead to

$$x_i(t_0 \leq \left( b_i^l - \sum_{\substack{j=1 \\ J\neq i}}^{n} a_{ij}^u \left(b_j^u / a_{jj}^l\right) \right) (1/a_{ll}^u) = x_i^l$$

$$\Rightarrow x_l(t) \geq x_i^l \ for\ t \geq t_0,\ i = 1, 2, \dots, n. \quad (1.7)$$

16. It is our aim to show that the mode together with other models stated in the foregoing together with haimovicis model is applicable to whole lot of consummative debit – credit systems in cosmology, astrophysics, finance, economics, sociology, anatomy computation. We will not however divide the system studied here into 9 components/ modules due to restrictions on space….

From the foregoing preparation we have the following:

LEMMA 1.1. *Assume that* $b_i, a_{ij} (i, j = 1, 2, \dots, n)$ *are scalar nonnegative almost periodic functions defined on* $\boldsymbol{R}$ *such that (1.2) and (1.3) hold. Then the set S defined by*

$$S = \{x = (x_1, x_2, \dots, x_n) \in \boldsymbol{R}_n | x_* \leq x_l \leq x^*; i = 1, 2, \dots, n\} \quad (1.8)$$

*Where*

$$x_* = \min_{1 \leq i \leq n} x_l^l; \qquad x_* = \min_{1 \leq i \leq n} x_i^u \quad (1.9)$$

*is invariant with respect to (1.1)*

**Appellation, cognomen, compellation, designation, denomination of the conducive, infelicitous and unfortutious** responses to environment change may be influenced by changes in available domain and milieu under consideration. An example is the following study: Authors describe the gravitational degrees of freedom of the Schwarzschild black hole by one free variable. Authors introduce an equation which we suggest to be the Schrödinger equation of the Schwarzschild black hole corresponding to this model. We solve the Schrödinger equation explicitly and obtain the mass spectrum of the black hole as such as it can be observed by an observer very far away and at rest relative to the black hole. Our equation implies that there is no singularity inside the Schwarzschild black hole, and that the black hole has a certain ground state in which its mass is non-zero. **Schrodinger equation of the Schwarzschild black hole Jarmo Mäkelä (Cambridge U E-Print: gr-qc/9602008 | PDF** Particles, axions, atoms, molecules, quarks, Majorana fermions, bosons et al in its confederation capacity and aggregation processes, type interactions and interactions between landscape dynamics. Following bears ample testimony, infallible observatory to the fact there might be changes in the quantum behaviour due to merger of black holes. Following example bears enough testimony, infallible observatory and impeccable demonstration and apodictic evidence as how some of the models could be applied to cosmological problems. This is attributable and ascribable to the fact that the classification criterion is based on the age.

**How gravity kills (e) Schrödinger's cat Theorists argue that warped space-time prevents (e) quantum superpositions of large-scale objects. Elizabeth Gibney 17 June 2015 Propositions are written in model able form and can easy be scheduled.**

| | |
|---|---|
| (155) | If the cat in Erwin Schrödinger's famous thought-experiment behaved according to quantum theory, it would be able to exist in multiple states at once: both dead and alive. |
| (156) | Physicists' common explanation for why we don't (e) see such quantum superpositions — in cats or any other aspect of the everyday world — is interference from (e) the environment. |
| (157) | As soon as a quantum object interacts with (e&eb) a stray particle or a passing field, it picks (e) just one state, collapsing (eb) into our classical, everyday view. |
| (158) | But even if physicists could completely isolate (eb) a large object in (eb) a quantum superposition, according to researchers at the University of Vienna, it would still collapse into (e&eb) one state — on Earth's surface, at least. |
| (159) | "Somewhere in interstellar space it could be that the cat has (e) a chance to preserve quantum coherence, but on Earth, or near any planet, there's (=) little hope of that," says Igor Pikovski. The reason, he asserts, is gravity. |
| (160) | Pikovski and his colleagues' idea, laid out in a paper published in Nature Physics on 15 June1, is at present only a mathematical argument. But experimenters hope to test whether gravity really does collapse (eb) quantum superpositions, says Hendrik Ulbricht, an experimental physicist at the University of Southampton, UK. "This is a cool, new idea, and I'm up for trying to see it in experiments," he says. Assembling the technology to do so, however, may take as long as a decade, he says. |
| (161) | Cinema-goers who saw the film Interstellar are already familiar with the basic principle behind the Vienna team's work. Einstein's theory of general relativity states that an extremely massive object causes clocks near it to (e) run more slowly because (e) its strong gravitational field stretches (e&eb) the fabric of space-time (which is why a character in the film aged only an hour near a black hole, while (e) seven years passed on Earth). On a subtler scale, a molecule placed nearer the Earth's surface experiences (e&eb) a slightly slower clock than one placed slightly further away. |
| (162) | Because of gravity's effect on (e&eb) space-time, Pikovski's team realised that variance in a molecule's position will also influence (e&eb) its internal energy — the vibrations of particles within (eb) the molecule, which evolve (eb) over time. |
| (163) | If a molecule were put in a quantum superposition of two places, the correlation between position and (e&eb) internal energy would soon cause (eb) the duality to 'decohere' to (e) the molecule taking just one path, they suggest. |
| (164) | "In most situations decoherence is due to (e) something external; here it's as though the internal jiggling is interacting with (e&eb) the motion of the molecule itself," adds Pikovski. |

**A practical limit**

| | |
|---|---|
| (165) | No one has yet seen this effect because other sources of decoherence — such as magnetic fields, thermal radiation and vibrations — are (=) typically much stronger and cause (eb) quantum systems to collapse long before gravity becomes an issue. But experimenters are keen to try. |
| (166) | Markus Arndt, an experimental physicist also at the University of Vienna, already tests whether quantum superpositions can be observed for (e) large, though not cat-sized, objects (see Fattening up Schrödinger's cat). |
| (167) | He sends large molecules through (e&eb) a matter-wave interferometer, a system that gives (eb) each molecule a choice of (e) two different paths. |
| (168) | In the classical view, a molecule travels down just one path; a quantum molecule effectively passes through (e&eb) both routes at once and interferes with (e&eb) itself to create (eb) a characteristic wave-like pattern (see image, below). |

| | |
|---|---|
| (169) | A quantum interference pattern built up from complex molecules known as phthalocyanines.(deleted due to spatial constraints) |
| (170) | A similar set up could be used to (e) test for gravity's ability to destroy (e) quantum behaviour: by comparing a vertical interferometer, in which the superposition should soon decohere due to (e) the stretching of time in one path versus (e&eb) the other, with a horizontal set up, where (e) the superposition could remain. |
| (171) | Arndt, who has tested the effect for (e) molecules of up to 810 atoms2, points out those large molecules, would be (=) good for testing the gravitational effect because (e) they contain (e) lots of particles contributing to (e) internal energy. |
| (172) | **But not only would the researchers have to further suppress the external environment to reduce (e) other decoherence effects, they would also need to (e) increase the separation of the two paths from micrometres to metres, or else use molecules perhaps 1 million times more massive. "It's certainly very challenging," says Arndt**. |
| (173) | If gravity's effect does limit (e) quantum behaviour on Earth, tests of quantum reality for (e) large objects may eventually have to (e) move into space, says Angelo Bassi, a physicist at the University of Trieste in Italy. |
| (174) | "But from a deep, fundamental point of view, this is nothing new," he says. A gravitational field is merely another environment to interact with, so invoking it does not (e) explain whether quantum behaviour might lead to (eb) classical reality if gravity's influence were mitigated — for example, by doing the experiment in (eb) gravity-free space. |
| (175) | Nor does the effect described by Pikovski and colleagues say anything about quantum gravity: a theory that would unite gravity and (e&eb) quantum mechanics in (eb) a single description, which many researchers are working towards. |
| (176) | "It's an interesting effect, but it's still quantum physics applied to (e&eb) classical general relativity. In that way it doesn't change our picture of the world," adds Bassi. Nature doi:10.1038/nature.2015.17773 |

Integrated factors these influences because they deal with ambience, mise en scene change or confederation dynamics, but rarely both. Time series of (e) space time suitability models with spatially explicit stochastic models to explore factors that influence the viability of type viability and confederation feasibility under stable and changing scenarios. Methodology could be applied to interrogation propositions and problems which might subsume one with the other notwithstanding that there is a space of singularities and there is space for completion of present from the past and future. An example is Time image movement from Deleuze. We clearly face a stiff challenge in our attempt to proffer Deleuzian treatments of photography as our critical project brushes up against the fact that Deleuze does not invest photography with (e&eb)the creative and re-creative potential that he associates with (e&eb) both cinema and modern painting. In short, while he highly values other visual artistic forms, he seemingly presents photographic texts as stagnate documents or tools that produce certainty, organize bodies and desires, and iterate hackneyed ideas. He even uses photography as something of a foil to demonstrate the innovation of the cinema and the originality of modern painters. In Francis Bacon: The Logic of Sensation (1981), he notes how "photography has taken over the illustrative and documentary role, so that modern painting no longer needs to fulfill this function, which still burdened earlier painters" (10). He later explains that "photographs are ways of seeing, and as such, they are illustrative and narrative reproductions or representations ... . They are what are seen, until finally one sees nothing else." He treats the photo as an instrument for reproducing representations of reality—a device that iterates images until they are ossified as established stories, icons, or even stagnant perceptions. Indeed, like Bergson before him, Deleuze compares photography to what he identifies as a reductionist model of perception: "what we see, what we perceive, are photographs" (74).[4] The photo is, in effect, an always-already passé

sensational experience that numbs our sensitivity to the ongoing vitality of life and provides us with an efficient understanding of the dynamism of our world. But it is compelling that Deleuze theorizes both the quotidian nature and the liberatory potential of photography; he treats it as a common, everyday visual experience that, despite its aesthetic and technological limitations, emancipates the modern painter to explore new artistic opportunities. This special issue of Rhizomes repeatedly provides artistic examples and critical strategies that invite us to see photography as replete with the power to unburden and emancipate that Deleuze associates with both film and painting. Indeed, one way to think of this project of reading Deleuze against Deleuze might be to ask: if photography can free painting and painters to pursue novel creative opportunities, can it likewise free itself, photographic practitioners, and even viewers to create and re-create anew? **Rhizomes » Issue 23 (2012) » Michael Kramp Unburdening Life, or the Deleuzian Potential of Photography Michael Kramp Lehigh University** dmk209@lehigh.edu

**Existence of an almost- periodic solution**

DEFINITION 2.1. A function $f(t,x)$,where $f$ is an $m$- vector, $t$ is a real scalar and $x$ is an $n$- vector, is said to be almost periodic in $t$ uniformly with respect to $x \in X \subset R_n$, $if\ f(t,x)$ is continuous in $t \in R\ and\ x \in X$, and if for any $\varepsilon > 0$, it is possible to find a constant $l(\varepsilon) > 0$ such that in any interval of length $l(\varepsilon)$ there exists a $\tau$ such that the inequality

$$\left||f\ (t+\ \tau,x) - f(t,x)|\right| = \sum_{l=1}^{m} |f_l\ (t+\ \tau,x) = f_l(t,x)|$$

$$< \varepsilon \qquad (2.1)$$

is satisfied for all $t \in (-\infty,\infty), x \in X$. The number $\tau$ is called an $\varepsilon -$ translation number of $f(t,x)$.

Let $\{\lambda_j\}$ denote the set of all real numbers such that

$$\lim_{T\to\infty} T^{-1} \int_0^T f(t,x)\exp(-i\,\lambda_j t)dt \not\equiv 0 \quad for\ x \in X; i = \sqrt{-1}. \qquad (2.2)$$

It is known that when the set X is separable, the set of numbers $\{\lambda_j\}$ in (2.2) is countable. The set $\{\sum_1^N n_j\ \lambda_j\}$ for all integers N and integers $n_j$ is called the module of $f(t,x)$.

The following result concerned with module containment will be used in the proof of our existence theorem below.

**MODULE CONTAINMENT THEOREM** (Yoshizawa [7], page 18). *Let $f(t,x)$ and $g(t,x)$ be almost periodic in t uniformly for $x \in D \subset \boldsymbol{R}_n$. If for any compact set $S \subset D$ and for any sequence of real numbers $\{\tau_k\}$ having a limit as $k \to \infty$ (the limit being finite or$\pm\infty$) for which $\{f(t+\tau_k,x)\}$ is uniformly convergent on $\boldsymbol{R} \times$*

$S, \{g(t+,x)\}$ *is also uniformly convergent on* $\boldsymbol{R} \times S$ *then the module of* $g(t+,x)$ *is contained in the module of* $f(t+,x)$.

We can now formulate our existence theorem.

**THEORM 2.1** *Assume that the almost-periodic parameters* $b_1, a_{1J}$ $(i,j = 1,2,\dots,n)$ *of (1.1) satisfy (1.2)-(1.3). Furthermore assume that there exists a positive number* $\mu$ *such that*

$$\inf_{1 \in \boldsymbol{R}} a_{11}(t) > \sum_{\substack{J=1 \\ J \neq 1}}^{n} (\sup_{t \in \boldsymbol{R}} a_{Jl}(t)) + \mu, \quad i = 1,2,\dots,n. \tag{2.3}$$

*Then the system (1.1) has a strictly positive (component wise) almost-periodic solution, say* $w(t) = w_1(t), \dots, w_n(t), t \in \boldsymbol{R}$, *whose module is contained in that of*

$$f(t,x) = \{f_1(t,x_1,\dots,x_n), f_2(t,x_1,\dots,x_n), \dots, f_n(t,x_1,\dots,x_n)\}$$

*Where*

$$f_l(t,x_1,\dots,x_n) = x_l \left\{ b_l(t) - \sum_{j=1}^{n} a_{lj}(t) x_J \right\},$$

$$i = 1,2,\dots n; (x_1,x_2,\dots,x_n) \in S. \tag{2.4}$$

PROOF. The set $S$ defined by (1.8)-(1.9) is compact in $\boldsymbol{R}_n$. Let $\{\tau_h\}$ be any sequence of real numbers such that $\tau_k \to \infty$ as $k \to \infty$ and

$$\sum_{l=1}^{n} |f_i(t+\tau_k, x_1, \dots, x_n) - f_l(t, x_1, \dots, x_n)| \to 0 \quad \text{as } k \to \infty \tag{2.5}$$

Uniformly in $(x_1, \dots, x_n) \in S$ for $t \in (-\infty, \infty)$. for any real number $\beta$let $k_0(\beta)$ be the smallest value of $k$ such that $\tau k_0 + \beta \geq t_0$. Since the invariance of $S$ implies that any solution of (1.1)

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), x_2(t+\tau_k, t_0, x_2(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\} \in S$$

For all $t \geq \beta$ and $k \geq k_0(\beta)$ whenever $\{x_1(t_0), \dots, n_n(t_0)\} \in S$, our first task in the proof is to show that the sequence

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Converges to a continuous bounded function say $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined for $t \in [\beta, \infty)$ such that $w_1(t) \in [x_*, x^*], t \in [\beta, \infty), i = 1, 2, \dots, n$ and the convergence is uniform on all compact subsets of $[\beta, \infty)$. Since $x_*$ and $x^*$ are independence of $\beta$, it is enough to show that

$$\{x_1(t+\tau_k, t_0, x_1(t_0)), \dots, x_n(t+\tau_k, t_0, x_n(t_0))\}, \qquad k \geq k_0(\beta),$$

Is a Cauchy sequence on compact subsets of $[\beta, \infty)$.

Let $U$ be any compact subset of $\beta, \infty)$., and let be any arbitrary positive number. Choose an integer $n_0(\varepsilon, \beta) \geq k_0(\beta)$ so large that for $m \geq k \geq n_0(\varepsilon, \beta)$

and all $t \in (-\infty, \infty)$

$$\sum_{i=1}^{n} |b_1(t+\tau_m) - b_1(t+\tau_k)| < (\mu\varepsilon x_*)/(4x^*), \tag{2.6}$$

$$x^* \sum_{i=1}^{n}\sum_{j=1}^{n} |a_{lj}(t+\tau_m) - a_{lj}(t+\tau_k)| < (\mu\varepsilon x_*)/(4x^*) \tag{2.7}$$

$$(2nx^*/x_*)\exp[-\mu x_*(\beta + \tau_k - t_0)] < (\varepsilon/2x^*)$$

(2.8)

Denote $x_i(t, t_0, x_1, (t_0))$ $by$ $x_1(t)$ brevity, and consider a function $v(s)$ defined by

$$v(s) = v\left(x(s), x(s+\tau_m+\tau_k)\right)$$

$$= \sum_{i=1}^{n} |\log x_1, (s) - \log x_1(s+\tau_m+\tau_k)|, \qquad x \geq t_0, \tag{2.9}$$

Where $x(t) = \{x_1(t), \dots, x_n(t)\}$ denotes a solution of (1.1) with $\{x_1, (t_0), \dots, x_n(t)\} \in S$. the invariance of S for (1.1), together with the elementary mean value theorem of differential calculus, implies

$$(1/x^*)\sum_{l=1}^{n}|x_l(s)-x_1(s+\tau_m+\tau_k)|\ \upsilon(s)$$

$$(1/x^*)\sum_{i=1}^{n}|x_l(s)-x_1(s+\tau_m+\tau_k)|. \tag{2.10}$$

Calculating the right derivative $D^+\upsilon$ of $\upsilon$ along the solutions of (1.1) and simplifying

$$D^+\upsilon(s)=\sum_{i=1}^{n}\{sign[x_i(s)-x_i(s+\tau_m-\tau_k)]\}$$

$$\left\{\frac{d[logx_i(s)]}{ds}-\frac{d[logx_i(s+\tau_m-\tau_k)]}{ds}\right\}$$

$$\sum_{l=1}^{n}|b_l(s)-b_l(s+\tau_m-\tau_k)|$$

$$+\sum_{l=1}^{n}(-a_{il}(s)|x_i(s)-x_i(s+\tau_m-\tau_k)|+\sum_{\substack{J=1\\ j\neq l}}^{n}a_{ij}(s)|x_j(s)-x_j(s+\tau_m-\tau_k)|)$$

$$+\sum_{i=1}^{n}\sum_{j=1}^{n}|a_{ij}(s)|a_{ij}(s+\tau_m-\tau_k)|\,x_j(s+\tau_m-\tau_k) \tag{2.11}$$

$$-\mu x_*\upsilon\big(x(s),x(s+\tau_m-\tau_k)\big)+\sum_{l=1}^{n}|b_l(s)-b_i(s+\tau_m-\tau_k)|$$

$$+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|. \tag{2.12}$$

Let us set $s-\tau_k=\sigma$; then we have from (2.6)-(2.7) that

$$\sum_{i=1}^{n}|b_l(s)-b_1(s+\tau_m-\tau_k)|+x^*\sum_{i=1}^{n}\sum_{J=1}^{n}|a_{ij}(s)-a_{ij}(s+\tau_m-\tau_k)|.$$

$$= \sum_{l=1}^{n} |b_l(\sigma + \tau_k) - b_i(\sigma + \tau_m)|$$

$$+x^* \sum_{i=1}^{n} \sum_{J=1}^{n} |a_{ij}(\sigma + \tau_k) - a_{ij}(\sigma + \tau_m)|$$

$$< (\mu\varepsilon x^*)/(2x^*). \quad (2.13)$$

Now an integration of (2.12) over $(t_0, t - \tau_k)$ together with (2.13) leads to

$$\upsilon\big(x(t+\tau_k), x(t+\tau_m)\big) \leq \upsilon\big(x(t_0), x(t_0 + \tau_m - \tau_k)\big) \exp[-\mu x_*(t + t_k - t_0)]$$

$$+[(\mu\varepsilon x_*)/(2x^*)]/\mu x_*)$$

$$\leq [(2x_n^*)/x_*]\exp[-\mu x_*(t + t_k - t_0)] + (\varepsilon/2x^*)$$

$$< (\varepsilon/x^*) \text{ On using (2.8), (2.10)} \quad (2.14)$$

For $m \geq n_0(\varepsilon, \beta)$ and for all $t \in \cup \subset [\beta, \infty]$, Thus we have from the definition of $\upsilon$ and (2.14) that

$$\sum_{l=1}^{n} |log x_i(t + \tau_k) - log x_l(t + \tau_m)| < (\varepsilon/x^*) \quad (2.15)$$

And hence using $x_* \leq x_i(t) \leq x^*, i = 1, 2, \dots, n; t \geq t_0$,

$$\sum_{i=1}^{n} |x_i(t + \tau_k) - x_i(t + \tau_m)| < \varepsilon \text{ for all } t \in \cup \subset [\beta, \infty], m \geq n_0(\varepsilon, \beta). \quad (2.16)$$

The existence of a function $w(t) = \{w_1(t), \dots, w_n(t)\}$ defined on $[\beta, \infty]$, such that $w_i(t) \in [x_*, x^*], i = 1, 2, \dots, n; t \in [\beta, \infty]$ follows from (2.16) and the fact that $x_i(t) \in [x_*, x^*], i = 1, 2, \dots, n;$ for $t \geq t_0$. Since $\beta$ is arbitrary, $w$ is in fact

Defined on $(-\infty, \infty)$ and hence we have

$x_i(t+\tau_k) \to w_l(t)$ as $k \to \infty$ uniformly on compact subsets of $(-\infty,\infty)$, $i = 1,2,\dots,n.$

(2.17)

Our next task is to show that $w: R \to S$ is differentiable and is a solution of (1.1). Since $x(t) = \{x_1(t),\dots,x_n(t)\}$ satisfies (1.1) we have $\sum_{i=1}^{n} |\frac{d}{dt} x_i(t+\tau_k) - \frac{d}{dt} x_i(t+\tau_m)|$

$$\sum_{i=1}^{n} |f_l\left(t+\tau_k, x_1(t+\tau_k),\dots,x_n\left(t+\tau_k\right)\right)$$

$$- f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)| \qquad (2.18)$$

$$+\sum_{l=1}^{n} |f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$- f_l\left(t+\tau_m, x_1(t+\tau_m),\dots,x_n(t+\tau_m)\right)|$$

Where $f_l\ (i = 1,2,\dots,n)$are as in (2.4). Since $x(t+\tau_k) = x_1(t+\tau_k),\dots, x_n(t+\tau_k) \in S$ for large $\tau_k$ for $t$ in each compact subset of $(-\infty,\infty)$ there exists an integer $n_1(\varepsilon) > 0$ such that if $m \text{ } k \text{ } n_1(\varepsilon)$, then for any $\varepsilon > 0$,

$$\sum_{i=1}^{n} |f_l\left(t+\tau_k, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$-f_1\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)| < (\varepsilon/2). \qquad (2.19)$$

Similarly since $x_l(t+\tau_k) \in S$ and $x_l(t+\tau_m) \in S$ for $i = 1,2,\dots,n$ there is an integer $n_2(\varepsilon) > 0$ such that for large enough $m, k$, such that $m \text{ } k \text{ } n_2(\varepsilon)$, we have

$$\sum_{l=1}^{n} |f_l\left(t+\tau_m, x_1(t+\tau_k),\dots,x_n(t+\tau_k)\right)$$

$$-f_1\left(t+\tau_m, x(t+\tau_m),\dots,x_n(t+\tau_m)\right)| < (\varepsilon/2) \qquad (2.20)$$

This equation will be very much useful in chemistry if used with the modules mentioned and written in concatenation with each other. It should be borne in mind that the applicability is modulewise and in no case concatenation should be marginalized. In all cases parameters are different which could be represented with different superscripts and subscripts like 1,23,…….etc., **Anal Bioanal Chem. 2004 Sep; 380(1):36-45. Epub 2004 Jul 29. Electronic tongue for pharmaceutical analytics: quantification of tastes and masking effects Legin A et al**

The organoleptic aspects of pharmaceutical formulations affect their acceptability to the patient and hence can have an important effect on concordance with treatment. Objective evaluation of these aspects, particularly the taste of the formulation and the drug substance it contains, is difficult. Whilst volunteer taste panels can be used to good effect their utility is limited, particularly during very early stage development when the toxicological profile of the active pharmaceutical ingredient (API) is yet to be established in detail. A potentiometric "electronic tongue" has been applied to analyse a variety of 41 individual substances and mixtures of particular interest for pharmaceutical research and development. The electronic tongue (ET) was capable of discriminating between substances with different taste modalities and could also distinguish different substances eliciting the same basic taste; the ET is promising in terms of quantifying the content of each substance and has an ability to detect nuances of the basic taste (e.g. lingering or short-lived). After calibration the electronic tongue was successfully applied to predicting bitterness strength of binary mixtures with a sweetener in terms of "apparent" or "perceived" quinine content. In order to render a formulation palatable it is often necessary to mask the (usually bitter) taste of the API by the addition of masking agents such as sweeteners and flavours. The ET proved capable of distinguishing between formulations with different levels of sweetener and/or flavour in a manner that was consistent with their masking efficiency as perceived by a small human taste panel. A suitably calibrated ET could have the benefit of providing the pharmaceutical formulator with reliable data concerning the taste of the product quickly and with a reduced need to ask volunteers to taste active pharmaceutical samples. Early development activities could be facilitated when human tasting is usually not possible in the absence of the required toxicological data.

Thus if $m \neq n_0(\varepsilon) = max\{n_1, n_2\}$, we have

$$\sum_{l=1}^{n} |\frac{d}{dt}\, x_1(t+\ \tau_k) - \frac{d}{dt} x_1\ (\ (t+\tau_m)| < \varepsilon \qquad (2.21)$$

Which shows that

$$\lim_{k\to\infty} \frac{d}{dt} x_1(t+\ \tau_k) \quad exists\ for\ i = 1, 2, \dots, n$$

Uniformly on all compact subsets of $(-\infty, \infty)$. Thus we have

$$\lim_{k\to\infty} \lim_{k\to\infty} \frac{x_i(t+\tau_k+h) - \ x_i(t+\tau_k)}{h}$$

$$= \lim_{h\to 0} \lim_{k\to 0} \frac{x_l(t+\tau_k+h) - \ x_l(t+\tau_k)}{h} \qquad (2.22)$$

$$= \lim_{h\to 0} \frac{w_1(t+h) - \ w_1(t)}{h}, \qquad i = 1, 2, \dots, n,$$

Showing that (since the left side of (2.22) exists) $w_i (i = 1, 2, \dots, n)$ is differentiable. Furthermore we have from (2.22) that

$$\frac{dw_l(t)}{dt} = \lim_{k\to\infty} [\, f_i\ \big(t+\tau_k, x_1(t+\tau_k), \dots, x_n\ (t+\tau_k)\big)$$

$$-f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big)$$

$$+ - f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big)]$$

$$= \ -f_l\big(t+\tau_k, w_1(t), \dots, w_n(t)\big); \qquad i = 1, 2, \dots, n, \qquad (2.23)$$

By the uniform almost periodicity of $f_1, f_2, \dots, f_n$ on $R \times S$ and hence $w(t) = \big(w_1(t), \dots, w_n(t)\big)$is a solution of (1.1).

In order to complete the proof, we have to show that $w(t) = \big(w_1(t), \dots, w_n(t)\big)$ is an almost- periodic function whose module is contained in that of $\{f_1(t, x_1, \dots, x_n), f_2(t, x_1, \dots, x_n), \dots, f_n(t, x_1, \dots, x_n)\}$ for $(x_1, x_2 \dots, x_n) \in S$. For this purpose it is in fact sufficient (by the module containment theorem) to show that for any sequence of real numbers $\{\tau_k\}$ for which

$$\{f_1(t + \tau_k, x_1, \dots, x_n), f_2(t, + \tau_k, x_1, \dots, x_n), \dots, f_n(t, + \tau_k, x_1, \dots, x_n)\}$$

Converges uniformly on R for all $(x_1, \dots, x_n) \in S$, the sequence $\{w_1(t, + \tau_k) \dots, w_n(t + \tau_k)\}$ converges uniformly on R where $\tau_k$ tends to a finite number or $\pm\infty$ as $k \to \infty$.

In the case where $\tau_k \to \tau \neq \pm\infty$ as $k \to \infty$, we have $w_l(t, + \tau_k) \to w_i(t + \tau), i = 1, 2, \dots, n$ uniformly on R as $k \to \infty$. Let us suppose $\tau_k \to \infty$ as $k \to \infty$. For any $\varepsilon > 0$ there exists an integer $n_0(\varepsilon)$ such that if $m$ [illegible] $n_0(\varepsilon)$ then

$$\sum_{l=1}^{n} |b_l(t + \tau_k) - b_l(t + \tau_m)| < (\varepsilon/4x^*)x_*, \qquad (2.24)$$

$$x^* \sum_{i=1}^{n} \sum_{j=1}^{n} |a_{ij}(t + \tau_k) - a_{ij}(t + \tau_m)| < \left(\frac{\varepsilon}{4x^*}\right)\mu x_* \qquad (2.25)$$

$$n(2x^*/x_*) \exp[\mu x_* \tau_k] < (\varepsilon/2x^*). \qquad (2.26)$$

Consider the function $v\big(w(s), w(s + \tau_m - \tau_k)\big)$ for $s \in [t, t + \tau_k]$ defined by

$$v\big(w(s), w(s + \tau_m - \tau_k)\big) = \sum_{l=1}^{n} |\log w_l(s) - \log w_l(s + \tau_m - \tau_k)|. \qquad (2.27)$$

Using the fact $w: \boldsymbol{R} \to S$ and $w$ is a solution of (1.1), calculate the right derivative $D^+v \ of \ v$;

$$D^+v\big(w(s), w(s + \tau_m - \tau_k)\big) = \sum_{i=1}^{n} \{\text{sign}[w_l(s) - w_i(s + \tau_m - \tau_k)]\}.$$

$$\left\{\frac{dw_i(s)}{ds} \cdot \frac{1}{w_l(s)} - \frac{dw_l(s + \tau_m - \tau_k)}{ds} \frac{1}{w_l(s + \tau_m - \tau_k)}\right\}$$

$$\leq -\mu \sum_{l=1}^{n} |w_l(s) - w_l(s+\tau_m-\tau_k)| + \sum_{l=1}^{n} |b_i(s) - b_l(s+\tau_m-\tau_k)|$$

$$+x^* \sum_{l=1}^{n}\sum_{J=1}^{n} |a_{lJ}(s) - a_{lJ}(s+\tau_m-\tau_k)| \quad (2.28)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+ \sum_{i=1}^{n} |b_l(s) - b_l(s+\tau_m-\tau_k)| \quad (2.29)$$

$$+x^* \sum_{l=1}^{n}\sum_{j=1}^{n} |a_{lJ}(s) - a_{ij}(s+\tau_m-\tau_k)|.$$

If we set $s-\tau_k = \sigma$ then we have from (2.29),

$$D^+ v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+ \sum_{l=1}^{n} |b_l(\sigma+\tau_k) - b_l(s+\tau_m)|$$

$$+x^* \sum_{l=1}^{n}\sum_{j=1}^{n} |a_{ij}(\sigma+\tau_k) - a_{ij}(\sigma+\tau_m)| \quad (2.30)$$

$$\leq -\mu x_* v\big(w(s), w(s+\tau_m-\tau_k)\big)$$

$$+[(\varepsilon\mu x_*)/(4x^*)] + [(\varepsilon\mu x_*)/(4x^*)]$$

(2.31)

Integrating (2.31) over $(t, t+T_k)$

$$v(w(t+T_k)w(t+T_m))$$

Another MODEL where the model is directly applicable with the scheduled ones are the case of insurance and corporate finance. Insurance: Mathematics and Economics Volume 42, Issue 1, February 2008, Retrieval of Black–Scholes and generalized Erlang models by perturbed observations at a fixed time Dedicated to Professor André Dubey on the occasion of his 60th birthday Daniel Neuenschwander

S-stable laws on the real line (more generally on Hilbert spaces), associated with some non-linear transformations (so-called "shrinking operations"), were introduced in [Jurek, Z.J., 1977. Limit distributions for truncated random variables. In: Proc. 2nd Vilnius Conference on Probability and Statistics, June 28–July 3, 1977. In: Abstracts of Communications, vol. 3, pp. 95–96; Jurek, Z.J., 1979. Properties of s-stable distribution functions. Bull. Acad. Polon. Sci. Sér. Math. XXVII (1), 135–141; Jurek, Z.J., 1981. Limit distributions for sums of shrunken random variables. Dissertationes Math. vol. CLXXXV]. In [Jurek, Z.J., Neuenschwander, D., 1999. S-stable laws in insurance and finance and generalization to nilpotent Lie groups. J. Theoret. Probab. 12 (4), 1089–1107], the authors interpreted s-stable motions on the real line as limits of total amount of claims processes (up to a deterministic premium) of a portfolio of excess-of-loss reinsurance contracts and showed that they led to Erlang's model or to Brownian motion. In [Neuenschwander, D., 2000b. On option pricing in models driven by iterated integrals of Brownian motion. In: Mitt. SAV 2000, Heft 1, pp. 35–39], we considered stochastic integrals whose integrand and integrator are both independent Brownian motions, thus modelling a stochastic volatility; as a result we got an analogue of the Black–Scholes formula in this model, confirming a result of Hull and White [Hull, J., White, A., 1987. The pricing of options on assets with stochastic volatility. J. Finance XLII (2), 281–300]. In the present paper, we will look at a common generalization of these processes, namely s-stable motions on the real line perturbed by a stochastic integral whose integrand and integrator are both (not necessarily independent) s-stable motions. The main result will be that if we can observe the distribution of such so-perturbed s-stable motions (together with the values of the perturbing processes) at time t=1, then we can identify the whole model (including the perturbation) among all models with Lévy processes perturbed by an iterated stochastic integral of two Lévy processes (in the Gaussian case) resp. among all models with a compound Poisson process with drift perturbed by an iterated stochastic integral of two compound Poisson processes (in the completely non-Gaussian case if the perturbing processes have no drift) without knowing anything about the history or about its distribution during $0 \le t < 1$. This applies, e.g., to a situation

where several assets obey the same model and one can estimate the distribution at time one by looking at the values of all these assets at time t=1. Interestingly enough, it will be convenient to treat the whole matter in the algebraic framework of the so-called Heisenberg group. This is a concept coming in fact from quantum mechanics and is in a certain sense the simplest non-commutative Lie group.

$\leq e^{-\mu x^* T_k} v\big(w(t), w(t+T_m-T_k)\big) + (\varepsilon/2x^*)$

$\leq e^{-\mu x^* T_k} \sum_{i=1}^{n} |\log\ w_i(t) - \log\ w_i(t+T_m-T_k)| + (1/x_*) + (\varepsilon/2x^*)$

$\leq n(2x^*/x_*)\exp[-\mu x_* T_k] + (\varepsilon/2x^*)$

$<(\varepsilon/x^*)$.

Since

$$v\big(w(t+T_k)w(t+T_m)\big) \leq \sum_{i=1}^{n} |\ w_1(t+T_k) - w_1(t+T_m))/x^*)$$

We have from(2.33)that

$$\sum_{i=1}^{n} |\ w_1(t+T_k) - w_1(t+T_m)| < \varepsilon\ for\ m \geq k \geq n_0(\varepsilon)$$

Which show that the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\}$converges uniformly on **R**. By a similar argument the above conclusion is also valid when $T_k \to -\infty\ as\ k \to \infty$. Thus the uniform convergence of the sequence $\{w_1(t+T_k),\ldots., w_n(t+T_k)\}\ as\ k \to \infty\ on$ **R** follows and this completes the proof of the theorem.

**COROLLARY.** Suppose $b_i a_{ij}$ (i, j = 1.2, … , n)in (1.1)are periodic with a common period say $T > 0$. *Then if* (1.2), (1.3)*and* (2.3)*hold the system* (1.1)*will have a periodic solution* $w(t) = \{w,(t), \ldots\ldots, w_n(t)\}$*of period* $T$ *such that*

$$w_1(t+T) = w_1(t), \qquad t \in \boldsymbol{R}: i = 1,2, \ldots, n.$$

**PROOF.** Proof follows from that of Theorem 2.1 if we Choose $T_k = KT(K = 0,1,2, \ldots)$*and note that*

$$x_1(t+KT) - w_1(t) \to 0\ as\ k \to \infty; i = 1,2, \ldots, n; t \in \boldsymbol{R},$$

$$x_i(t+(K+1)t) - w_1(t) \to 0 \ \ as \ \boldsymbol{K} \to \infty; i = 1,2,\dots,n; t \in \boldsymbol{R},$$

$$x_i(t+T+KT) - w_1(t+T) \to 0 \ \ as \ k \to \infty; i = 1,2,\dots,n; t \in \boldsymbol{R},$$

Implying that

$$w_1(t+T) = w_i(t); \quad i = 1,2,\dots,n; t \in \boldsymbol{R},$$

## 3. Global asymptotic stability

By definition we say that an almost-periodic solution say $w(t) = \{w_1(t),\dots,w_n(t)\} of (1.1) is\ globally\ asymptotically$ stable (or attractive) if and only if every other solution x(t)={ $x_1(t),\dots,x_n(t)$}with $x_1(o) > 0, i = 1.2,\dots,n$ is defined for all t□0 and satisfies

$$\lim_{t\to\infty} |w_1(t) - x_1(t)| = 0; \quad i = 1,2,\dots,n.$$

A consequence of such a global asymptotic stability is that there cannot be another strictly positive almost-periodic solution of (1.1)

**THEOREM 3.1** Suppose the almost- Periodic coefficients $b_1$, $a_{ij}$ (i, j = 1,2, … , n)

Satisfy (1.2)and (1.3)and the following

$$\inf_{t\in R} a_{ij}(t) > \left\{ \sum_{\substack{J=1 \\ J\neq 1}}^{n} \sup a_{ij}(t) \right\} + \mu; \qquad i = 1,2,\dots,n$$

For some positive number $\mu\ Then$ (1.1)has an almost-periodic solution (with strictly positive components) which is globally asymptotically stable.

We will omit the proof since the proof is identical with that of Theorem 3.1 of [3]

Under condition (3.2) the almost –periodic solution $w(t) = \{w_1(t), \ldots . w_n(t)\}$ of (1.1)enjoys a somewhat strong from of stability even if the coefficients of (1.1)are subjected to a class of perturbations We note (see Yoshizawa [7]page17)that if a function b:R→ $R$ is almost periodic then H(b),the "hull of b" is defined as follows:

H(b)= g|g:**R**→ $\boldsymbol{R}$ $and\ g(t)\lim_{k\to\infty} b(t+T_k) for$ some sequence

$T_k$ of real numbers the limit being uniform in $t \in R$.

By Theorem 2.2 of Yoshizawa [7, page 10], members of the hull of almost periodic functions are themselves almost periodic.

A system of the from

$$\frac{dy_1(t)}{dt} = y_1(t)\left\{B_i(t) - \sum_{J-1}^{n} A_{ij}\ (t)y_1(t)\right\},$$

*i=1,2,…,n; t>*$t_0 \in R$,

Where

$$B_i \in H(b_i); (A_{ij}), \qquad \text{i, j} = 1,2, \ldots, \text{n},$$

is known as a system from the hull of (1.1) As a consequence of the definition of the hull in (3.3)and (3.5)we shall have

$$\inf_{t\in R} B_i(t) = \text{b}_i^{\,l}; \quad \sup_{t\in R} B_i(t) = \text{b}_i^{\,l};$$

$$\inf_{t\in R} A_{ij}(t) = \text{b}_i^{\,l}; \quad \sup_{t\in R} A_{ij}(t) = \text{a}_{ij}^{\,u};$$

and hence the coefficients of(3.4)- (3.5)satisfy the same conditions as (1.2),(1.3)and (3.2)It will then follow from theorems(2.1)and (3.1)that solutions of (3.4)-(3.5)with $y_1(t_0) > 0, i = 1,2, \ldots, n\ will\ have$ the following property:

$$y_1(t) > 0 \ \ for\ t \geq t_0; i = 1,2,\dots,n,$$

$$\sum_{i=1}^{n} y_1(t) remains\ bounded\ for\ t \geq t_0$$

**DEFINITION** A Solution $x(t) = \{x_1(t), \dots, x_n(t)\}, t \in t_0 \in \boldsymbol{R}$,of (1.1)is said to be stable under perturbations from the hull of (1.1)if the following holds :let $y(t) = \{(y_1(t), \dots, y_n(t)\}$ be any solution of (3.4)-(3.5)such that $y_1(t_0) > 0$, i=1,2,…,n,and for some positive constants $\beta_1, \beta_2 \beta_1, \leq y_1(t) \leq \beta_2, t \geq t_0, \ \ i = 1,2,\dots,n$ Then given any $\varepsilon > 0$,there exists a constant $\delta > 0$ such that

$$\sum_{j=1}^{n} |x_1(t_0) - y_1(t_0)| < \delta$$

$$\sum_{i=1}^{n} \sup_{t \in R} |b_1(1) - B_1(t)| < \delta \quad = \sum_{i=1}^{n} |x_1(t) - y_i(t)| < \varepsilon$$

$$\sum_{i=1}^{n} \sum_{j=1}^{n} | \quad \sup_{t \in R} |a_{ij}\ (t) - A_{ij}(t)| < \delta$$

For $t \geq t_0$.

We can now formulate the following result which shows that all solution of (1.1)are stable under perturbations from the hull of (1.1)

**THEOREM 3.2**. Assume that the conditions of Theorem (3.1) hold. Then every strictly positive solution of (1.1) ( including its unique almost-periodic solution) is stable under perturbations from the hull of (1.1).

PROOF. Let $x(t) = \{x_1(t), \dots, x_n(t)\}$ and $y(t) = \{y_1(t), \dots, y_n(t)\}$ be any two strictly positive solutions of (1.1) and (3.4) – (3.5) respectively for $t \geq t_0 \in R$. Let $\beta_1, \beta_2$ be constants such that

$\beta_1 \leq x_i(t) \leq \beta_2$ ,
$\beta_1 \leq y_i(t) \leq \beta_2$ , $i = 1,2,\dots,n; t \geq t_0$ (3.7)

Consider now a function $V(t) = V(x(t)y(t))$ defined by

$$V(x(t)y(t)) = \sum_{i=1}^{n} |\log x_1(t) - \log, y_i(t)|, \quad t \geq t_0 \in \text{R}. \tag{3.8}$$

It is easy to see that

$$\frac{1}{\beta_2} \sum_{i=1}^{n} |x_i(t) - y_i(t)| \leq V(x(t)y(t))$$

$$\leq \frac{1}{\beta_1} \sum_{i=1}^{n} |x_i(t) - y_i(t)|, \; t \geq t_0 \in \text{R}. \tag{3.9}$$

Calculating the right derivative $D^+V$ of $V$ and using (1.1) and (3.4) we derive (after simplification) that

$$D^+V\ (x(t), y(t)) \leq -\mu \sum_{i=1}^{n} |x_i(t) - y_i(t)| + \sum_{i=1}^{n} |b_i(t) - B_i(t)|$$

$$+\beta_2 \sum_{j=1}^{n} \sum_{i=1}^{n} |a_{ij}(t) - A_{ij}(t)|; \quad t \geq t_0,$$

$$\leq -\mu\, \beta_1\, V\,(x(t), y(t)) + \sum_{i=1}^{n} |b_i(t) - B_i(t)|$$

$$+ \beta_2 \sum_{i=1}^{n} \sum_{j=1}^{n} |a_{ij}(t) - A_{ij}(t)|; \quad t \geq t_0, \tag{3.10}$$

An integration of (3.9) over $[t_0, t]$ with an application of (3.8) leads to

$$\frac{1}{\beta_2}\sum_{i=1}^{n}|x_i(t)-y_i(t)| \leq V(x(t),y(t)) \leq V(x(t_0),y(t_0)) + \left[\sum_{i=1}^{n}\sup_{t\in R}|b_i(t)-B_i(t)| + \beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}|a_{ij}(t)-A_{ij}(t)|\right]\frac{1}{\mu\beta_1},$$

$$t \geq t_0,$$

And hence

$$\sum_{i=1}^{n}|x_1(t)-y_1(t)| \leq \left(\frac{\beta_2}{\beta_1}\right)\sum_{i=1}^{n}|x_1(t_0)-y_1(t_0)| + \frac{\beta_2}{\mu\beta_2}\left[\sum_{i=1}^{n}\sup_{t\in R}|b_i(t)-B_i(t)|\right.$$

The construction reason is because is criticized in a number of usage guides: The reason for the long delays was because the costs greatly exceeded the original estimates. One objection to this construction is based on its redundancy: the word because (literally, by cause) contains within it the meaning of reason; thus saying the reason is because is like saying “The cause is by cause,” which would never be said. A second objection is based on the claim that because can introduce only adverbial clauses and that reason is requires completion by a noun clause. Critics would substitute that for because in the offending construction: The reason for the long delays in completing the project was that the costs.

$$\left.+\beta_2\sum_{i=1}^{n}\sum_{i=1}^{n}\sup_{t\in R}|a_{ij}(t)-A_{ij}(t)|\right]. \qquad (3.11)$$

Now for any $\varepsilon > 0$ if one chooses a $\delta > 0$ such that

$$\delta < \min\left\{\frac{\varepsilon}{3}\frac{\beta_1}{\beta_2}, \frac{\varepsilon}{3}\frac{\beta_1}{\beta^2}\mu\right\} \qquad (3.12)$$

Then (3..6) will be satisfied and this completes the proof

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# . Jia Li's Prescientiial Dynamics And Contential Staticity Of Accentuation Attrition Models Applied To Cosmological And Astrophysical And Quantum Mechanical Engineering Problems

**Abstract**: **Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics ET AL ). We shall specifically delineate and disseminate the information and particles and confederation of the financial systems to which the model is applied and in concomitant elucidation of the Haimovici's model we apply this to each and every module. Due to spatial constraints we have not done any separate sets as has been done in the foregoing. It is to be repeated that the model is to be applied for each and every module. In amplification this is true for all the (confederation of corresponding addition is made to the credit –debit system of atoms, molecules, qubits, cosmological interactions, financial transaction, economic exchange confederation as the case may be( Note the author deals for predator – population.) In amplification the model is applicable to various assortment of systems and their perspicuous characterstics, sophisticated seasoning, constituent structure, concept formulation, beat cadence, interiority of perfect causes, incompatibility of conjunctions', resonantial resumptions, hypothetical propositions, logical compatibilities, on causal correspondences, predicational inherence, emphasis related phenomenological methodologies, Partcipational observation, background actuality, addressed aestheticism, polished perception, refined proficiency, existential worldliness, existential strife, predicational anteriorities, apophthegmatic axiom, brocard- gnome, implantation and inculcation, pronouncement and proposition, declaration dogma, character consonance, shift and stratagem, ontological consonance, primordial exactitude, phenomenological correlates, accolytish representation, atrophied asseveration, suppositional surmise, theorization anamensial alienisms, anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orientationalities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state. Characterstics of such systems are taken in to consideration in the complete consolidation and consummation of the systems. Predator prey model literally form bastion, pillar, post, stylobate, sentinel towards the end of dismasting many of the asymptotic and solutional behaviour in finance, economic, and as we shall see in cosmology, quantum computation and other fields.**

## MODEL FORMULATION

Let $p_1(t, a_1)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_1}{dt} = -b_1 Q_1 + g_1(g_1)Q_1$$

Where $b$ is the natural death rate in the absence of credit types and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_1(0) = 0, \quad g_1(P_1) > 0, \ P_1 \ 0, \qquad \lim_{P_1 \to \infty} g_1(P_1) = (K_1)_1 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_1 + (\rho_a)_1 + D_1(a_1, P_1, Q_1)\rho_1 + 0, \qquad 0 < a_1 \ \infty,$$

$$\rho_1(t, 0) = \int_0^{\infty} n_1(a_1, P_1, Q_1)\rho_1(t_1, a_1)da_1$$

Where $D_1$ is the mortality function and $n_1$ is the fertility function.

Put $_1$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_1}{dt}\right]_1 = x_1\left[\frac{dQ_2}{dt}\right]_1 - x_2\left[\frac{dQ_1}{dt}\right]_1$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_1 - x_4\left[\frac{dQ_2}{dt}\right]_1$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_1 - x_6\left[\frac{dQ_3}{dt}\right]_1$$

Where $x_1, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_1$ and $g_1(P_1)Q_1$

Now when you add the term $(b_1Q_1)$ with $g_1(P_1)(Q_1)$ we are infect see is that $(b_1Q_1)$ is credited incrementally part of $g_1(P_1)(Q_1)$ or instant nearly, when subtracted from $\frac{dQ_1}{dt}$, it means that $g_1(P_1)(Q_1)$ is slowly eaten by $\frac{dQ_1}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_1 = \mu_1(a_1) + (f_1)_1(P_1)_1 + (f_2)_1(Q)_1 \quad (2)$$

$$n_1 = \beta_1(a_1)h_1(Q_1), \quad (3)$$

Where $\mu_1$the age- dependent debit rate whit is $\mu_1(a_1)$ 0 $\forall a > 0$; $(f_1)_1$ is the density- dependent death rate with $f_1(0) = 0$, $(f_1^1)_1(P_1)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_1)_1(P_1) = (k_1)_1$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_1(0) = 0, (f_2^1)_1(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_1(Q_1) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^1(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_1 + (\rho_a)_1 = -\mu_1(a_1) + (f_1)_1P_1) + \big(f_2)_1(Q_1)\big)(\rho_1)_1 \quad (4)$$

$$\rho_1(t_1, 0) = h_1(Q_1)\int_0^\infty \beta_1(a_1)\rho_1(t, a_1)da, \quad (5)$$

$$\frac{dQ_1}{dt}$$

$$= -b_1Q_1 + g_1(P_1)Q_1 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_1$ and $Q_1$ of (4-6) satisfy

$$0 \leq \rho_1(t, a_1) \leq \rho_1^*(t, a_1),$$

$$0 \leq Q_1(t) \leq Q_1^*(t),$$

Where $\rho_1^*(t, a_1)$ and $Q_1^*(t)$, are solutions of

$$(\rho_t^*)_1 + (\rho_a^*)_1 = -(\mu_1(a_1) + (f_1)_1(P_1^*) + (f_2)_1 Q_1^*))\, \rho_1^*, \quad (7)$$

$$\rho_1^*(t, 0) = \int_0^\infty \beta_1(a_1)\, \rho_1^*(t, a_1) da_1 \quad (8)$$

$$\frac{dQ_1^*}{dt} = -b_1 Q_1^* + g_1(P_1^*) Q_1^*, \quad (9)$$

With $P_1^*(t) = \int_0^\infty \rho_1^*(t, a_1) da_1$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_1^*(t)}{dt} = P_1^*(t)\big(r_1 - (f_1)(P_1^*) - (f_2)_1(Q^*)\big), \quad (10)$$

$$\frac{dQ_1^*(t)}{dt} = Q_1^*(t)(-b_1 + g_1(P_1^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_1(a) e^{-\int_0^a \mu_1(\tau) dt} e^{-ra}\, da_1 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_1^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(t)dt}\, da_1 > 1, \qquad (H1)$$

Which guarantees $r_1 > 0$.

*Remark.* If $\beta_1(a) = [\beta_0]_1 a_1 e^{-\alpha_1 a_1}, \mu_1(a_1) \equiv \mu_1$ then the solution of (12) is

$$r_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_1]_1(P_1) = 0, [f_2](Q_1) = \lambda_1 Q_1\ h_1(Q_1) = 1, \qquad and\ \ g_1(P_1) = c_1 P_1$$

$$[\rho_1]_1 + [\rho_a]_1 = -(\mu_1 + \lambda_1 Q_1)\rho_1$$

$$\rho_1(t,0) = \int_0^{\infty} [\beta_0]_1 a_1 e^{-\alpha_1 a_1} \rho_1(t, a_1)\, da_1$$

$$\frac{dQ_1}{dt} = -b_1 Q_1 + c_1 P_1 Q_1$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_1}{dt} = \omega_1 P_1 - \lambda_1 P_1 Q_1$$

$$\frac{dQ_1}{dt} = b_1 Q_1 - c_1 P_1 Q_1$$

Where $\omega_1 = \sqrt{\beta_1} - \alpha_1 - \mu_1$ THE EQUILIBRIUM $(\rho_0)_1(a_1), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_1(a_1), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_1 + (\rho_a)_1 = -\big(\mu_1(a_1) + (f_1)_1(P_1)\big)\rho,$$

$$\rho_1(t,0) = \int_0^\infty \beta_1(a_1)\rho_1(t,a_1)da_1 \qquad (13)$$

The solution of (13) has the form of

$$\rho_1(t,a_1) = Ce_1^{r(t-a_1)} - \int_0^a (\tau)d\tau - \int_0^1 (f_1)_1\big(P_1(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_1(t) = C_1 e^{r_1 t - \int_0^t (f_1)_1(P_1(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_1^{-1}(r)$ if $k_1 > r$.

This is equivalent to

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{(k_1)_1 a_1} da_1 < 1.$$

In fact, if we define

$$R_1(s_1) = \int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau)d\tau} e^{-s_1 a_1}\, da_1$$

$R(s)$ is a decreasing function. $R_1(r_1) = 1$ And $R_1(k_1)_1 < 1$ imply $r_1 < (k_1)_1$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_1(a_1), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{k_1 a_1} da_1 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_1 \equiv \int_0^\infty (\rho_0)_{11}(a_1)da_1$ *satisfies*

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau)d\tau} e^{-(f_1)_1(P_0)a_1} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_1(a_1), 0)$

Let $u_1(t, a_1) = \rho(t, a_1) - \rho_0(a_1)$ Then

$$N_1(t) \equiv \int_0^\infty u_1(t, a_1)da_1 = \int_0^\infty \rho_1(t, a_1)da_1 - \int_0^\infty (\rho_0)_1(a_1)da = P_1(t) - P_0)_1$$

And

$$(u_1)_1 + (u_a)_1 = (\rho_t)_1 + (\rho_a)_1 - (\rho_0^1)_1(a_1)$$

$$= -\mu_1(a_1) + (f_1)_1(P_1) + (f_2)_1(Q_1))(u_1 + (\rho_0)_1) + \mu_1(a_1)(\rho_0)_1 + (f_1)_1(P_0)_1(\rho_0)_1$$

$$= -\mu_1(a_1)u_1 - (f_1)_1(P_1) + (f_2)_1(Q_1))u_1 - (\rho_0)_1(f_1)_1(P_1) - (f_1)_1(P_0)_1$$

$$+ (f_2)_1(Q_1)), \qquad (15)$$

Where $(\rho_0^1)_1(a_1)$ denotes $d(\rho_0)_1(a_1)/da_1$

Expansion $(f_1)_1(P_1)$ and $(f_2)_1(Q_1)$ at $(P_0)_1$ and 0 yield respectively

$$(f_1)_1(P_1) = (f_1)_1((P_0)_1 + N_1)$$

$$= (f_1)_1(P_0)_1 + (f_1^1)_1(P_0)_1N_1 + o_1(N_1), \qquad (16)$$

$$(f_2)_1(Q_1)$$

$$= (f_2^1)_1(0)Q + o_1(Q_1) \qquad (17)$$

Write $f_1(P_0) = f_{10}, f_1'(P_0) = f_{10}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_1)_1 + (u_a)_1 \approx -(\mu_1(a_1) + (f_{10})_1)u - (f_{10}^1)_1(\rho_0)_1(N_1) - (f_{20}^1)_1(\rho_0)_1(Q_1)$$

$$u_1(t,0) = \rho_1(t,0)(\rho_0)_1(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t,a) + (\rho_0)_1(a_1)da_1 - (\rho_0)_1(0)$$

$$= h_1(Q_1)\int_0^\infty \beta_1(a_1)(u_1(t,a_1)da_1 + (h_1(Q_1) - 1)(\rho_0)_1(0) \qquad (18)$$

Expanding

$$h_1(Q_1) = 1 + (h^1)_1(0)Q_1 + 0(Q_1) \equiv 1 + h_0'Q_1 + 0Q_1$$

And substituting it into (18), we arrive at

$$u_1(t,0) \approx \int_0^\infty \beta_1(a_1)u_1(t,a_1)da_1 + (h_{a_1}')(\rho_0)_1Q_1$$

The linearization is the following:

$$[u_1] + [u_1] = -(\mu_1(a_1) + [t_{10}]_1u_1 - [\alpha_1]_1(a_1)N_1 - [\alpha_2]_1(a_1)Q_1$$

$$u_1(t,0) = \int_0^\infty \beta_1(a_1)u_1(t,a_1)da_1 + [\xi_1]Q_1 \qquad (19)$$

$$\frac{dQ_1}{dt} = -b_1Q_1$$

where

$$[\alpha_1]_1(a_1) = [t_1^1 a]_1[p, o]_1(a_1)d2(a_1) = [f_{20}^1]_1(\rho\, o\,)_1(a_1)\xi = [\mathrm{h\ o}^1]_1[\rho \mathrm{o}]_1(\mathrm{o}) \text{ and } \mathrm{b}_1 = \mathrm{b}_1 - g_1(P_1)$$

If $g_1(P_0) > b$, clearly, $(\rho_0)_1(a_1), 0)$ is unstable and the predator population is persistent. Assume $g_1(P_0)_1 < b$.

Then, $Q_1(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_1 + u_a = -(\mu_1(a_1) + (f_{10})_1)u_1 - (\alpha_1)_1(a_1)N_1$$

$$= -(\mu_1(a_1) + f_{10})u_1 - \alpha_1(a_1)\int_0^\infty u_1(t, a_1)da_1$$

$$u_1(t, 0) = \int_0^\infty \beta_1\,(a_1)u_1(t, a_1)da_1$$

Using the method developed by Cushing [2, 3], we suppose

$$u_1(t, a_1) = \tilde{\mathrm{u}}_1(a_1)^{e^{\gamma(t-a_1)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_1}\, da_1$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_1)$ gives,

$$\tilde{\mathrm{u}}_1(a_1) = e^{-\int_0^a \mu_1(\tau_1)d\tau - (f_{10})_1 a}\left(\tilde{\mathrm{u}}_1(0) - w\int_0^a e^{\int_0^\tau \mu_1(s_1)ds + (f_{10})}\tau_1\,(\alpha_1)_1(\tau_1)e^{\gamma_1\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_1(a_1)e^{-\gamma_1 a_1 - (f_{10})_1 a_1 - \int_0^a \mu_1(\tau_1)d\tau}\right)\tilde{\mathrm{u}}_1(0)$$

$$+w\int_0^{\infty}\beta_1(a_1)e^{-\gamma_1 a_1-(f_{10})_1 a_1-\int_0^a \mu_1(\tau_1)d\tau}\int_0^a(\alpha_1)_1(\tau_1)e^{-\gamma_1\tau_1}e^{(f_{10})_1\tau_1+\int_0^{\tau}\mu_1(s_1)ds}d\tau_1\,da_1$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty}e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}da\right)u(0)$$

$$+\left(-1-\int_0^{\infty}e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}\int_0^a e^{f_{10}\tau+\int_0^{\tau}\mu(s)ds+y\tau}\alpha_1(\tau)d\tau\,da\right)w=0.\ (23)$$

Define $Y(\gamma,a)=e^{-ya-f_{10}a-\int_0^a\mu(\tau)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta\,(a)Y(\gamma,a)da\right)u(0)$$

$$+w\int_0^{\infty}\beta(a)Y\,(\gamma,a)\int_0^a\alpha_1(\tau)Y^1(\gamma,\tau)d\tau\,da=0,(24)$$

$$\left(\int_0^{\infty}Y(\gamma,a)da\right)u(0)$$

$$+\left(-1-\int_0^{\infty}Y(\gamma,a)\int_0^a\alpha_1(\tau)Y^{-1}(\gamma,\tau)d\tau\,da\right)w=0.$$

Equation (20) has a solution of the form of $u(t,a)=u(a)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $u(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I(y,a)=\int_0^a\rho_0(\tau)Y^{-1}\,(\gamma,\tau)d\tau$. We obtain the characteristic equation

$$f_{10}^{l}\left(\int_0^{\infty} Y(\gamma,a)da\right)\left(\int_0^{\infty}\beta(a)Y(\gamma,a)I(\gamma,a)da\right)$$

$$+\left(1-\int_0^{\infty}\beta(a)Y(\gamma,a)da\right)\left(1+f_{10}^{l}\int_0^{\infty}Y(\gamma,a)I(\gamma,a)da\right)=0. \quad (25)$$

Since

$$\rho_0(a)=\rho_0(0)e^{-\int_0^a \mu(\tau)d\tau - f_{10}a}$$

$$I(\gamma,a)=\int_0^a \rho_0(\tau)e^{\gamma\tau+f_{10}\tau+\int_0^{\tau}\mu(s)ds}d\tau$$

$$=\frac{\rho_0(0)}{\gamma}(e^{\gamma a}-1)$$

Let

$$E(a)=e^{-f_{10}a-\int_0^a \mu(\tau)d\tau}$$

Then $[Y_1]_1[\gamma_1]_1,a_1)=\ E_1(a_1)e^{-\gamma_1 a_1}$. by a simple calculation,

$$\int_0^{\infty}\beta_1(a_1)Y_1(\gamma_1,a_1)I(\gamma_1,a_1)da=\frac{[\rho_0]_1(0)}{\gamma}\left(1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}\ da\right);$$

$$\int_0^{\infty}Y_1(\gamma_1,a_1)I(\gamma_1,a_1)da=\frac{[\rho_0]_1(0)}{\gamma_1}\left(\int_0^{\infty}E_1(a_1)e^{-\gamma_1 a_1}\ da_1\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_1(0)}{\gamma_1}[f_{10}^{1}]\int_0^{\infty}E_1(a_1)e^{-\gamma_1 a_1}da_1\ (\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1)$$

$$+\left(1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right).$$

$$\left(1+\frac{(\rho_0)_1(0)}{\gamma_1}(f_{10}^{'})_1\left(\int_0^{\infty}E_1(a_1)da_1-\int_0^{\infty}E_1(a_1)e^{-\gamma_1 a_1}da_1\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\right)(1+\frac{(f_{10}^{'})_1}{\gamma_1}(\rho_0)_1(0)\int_0^{\infty}E_1(a_1)da_1)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_1(a_1)Y_1(\gamma_1 a_1)da_1=1-\int_0^{\infty}\beta_1(a_1)E_1(a_1)e^{-\gamma_1 a_1}da_1\neq 0,$$

Since $\int_0^{\infty}\beta_1(a_1)Y_1(\gamma_1 a_1)\int_0^{a}\alpha_1(\tau_1)Y^{-1}(\gamma_1,\tau_1)d\tau_1 da_1>0$. Therefore, the characteristic solution is

$$\gamma_1=-(f_{10}^{'})_1(\rho_0)_1\int_0^{\infty}E_1(a_1)da_1.$$

Since $f_1'(P_0)_1>0$, we obtain the following result

THEOREM 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_1),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_1),0)$ which is related to the extinction of debit summation as has been referred to in the foregoing and to The possible persistence of the prey

population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_1 + (\mu_1(a_1) + [f_1]_1(P_1^*) + [f_2]_1(Q_1^*))P_1^* = 0, \quad (26)$$

$$p_1^*(0) = h_1(Q_1^*)\int_0^\infty \beta_1(a_1)p_1^*(a_J)p_1^*(a_1)da_1 \quad (27)$$

$$g_1(P_1^*) - b_1 = 0, \quad (28)$$

Where $P_1^* = \int_0^\infty p_1^*(a_1)\, da_1$

Equation (28) yields $P^* = g^{-1}(b)$,

From (26),

$$p_1^*(a) = p_1^*(0)e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}$$

Substituting into (27), we need to solve

$$\left(1 - h_1(Q_1^*)\int_0^\infty \beta_1(a_1)e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(Q_1^*)a_1}da_1\right)p_1^*(0) = 0$$

For $Q_1^*$

Fix $p_1^*(0)$ and $P_1^*$ Let

$$H_1(z_1) = 1 - h_1(z_1)\int_0^\infty \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1}da_{,1}$$

And assume that

$$[k_3]_1 \int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(g_1^{-1}(b_1))a_1 - [k_2]_1 a_1}\, da_1 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty} [H_1(z_1)] > 0.$$

Clearly,

$$\frac{dH_1}{dz_1} = h_1(z)(\,[f_2^{'}]_1(z_1)\,) \int_0^{\infty} a_1\beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - f_1(P_1^*)a_1 - f_2(z_1)a_1} da_1$$

$-h^{'}(z_1) \int_0^{\infty} \beta_1(a_1)\, e^{\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1 - [f_2]_1(z_1)a_1}\ da_1$

By the assumptions on $\left[f_2^1\right]_1 and\ [h^1]^1 dH_1/dz > 0$. Hence, if

$$H_1(0) = 1 - \int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1(P_1^*)a_1} da_1 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_1(0) < 0$ follows if $g_1^-(b_1) < [f_1^{-1}]_1(r_1)$, where $r_1$ is the solution of (12).

Indeed, since

$$R_1(s_1) = \int_0^{\infty} \beta_1(a_1)\, e^{-\int_0^a \mu_1(\tau_1)d\tau_1 e^{-y\frac{a}{1}}} da_1$$

Is a decreasing function of $s_1$ when $P_1^* = [g^{-1}]_1(b_1) < [f_1^{-1}]_1(r_1)$, that is $[f_1]_1(P_1^*) < r$, $R\big(f_1(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_1(g^{-1}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1)d\tau_1 - [f_1]_1\left(g_1^{-1}(b)\right)a_1} > 1. \qquad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_1^*(a_1), Q_1^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_1^*(a_1), Q_1^*)$ is more complicated. We again consider perturbations $m_1(t_1 a_1) = p_1(t_1, a_1) - p_1^*(a_1)$ and $W_1(t) = Q_1(t) - Q^*$.Then

$$M_1(t) \equiv \int_0^\infty m_1(t, a_1)da_1 = P_1(t) - P_1^*,$$

$$[m_1]_1 + [m_a]_1 = [p_1]_1 + [p_a]_1 - \left[p^*\right]_1(a_1) \approx -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1$$

$$-[f_{1*}^1]_1 p_1^*(a_1)M - [f_{2*}^1]_1 p_1^*(\mathrm{a})W_1,$$

$$m_1(t,0)p_1(t,0) - p_1^*(0) \approx h_* \int_0^\infty \beta_1(a_1)m_1(t,a_1)\,da + \frac{[h_*^1]_1}{[h_*]_1}\, p_1^*(0)W_1$$

$$\frac{dW_1}{dt} = \frac{dQ_1}{dt} \approx [g_*^1]_1 Q_1^* M_{,1}$$

Where $[f_{1*}]_1$ denoting $[f_1(P^*)]_1$, $[f_{1*}]_1$ denoting $[f_2(P^*)]_1$, $\left[f_{1*}^{'}\right]_1$ denoting $\left[f^{'}{}_1(P^*)\right]_1 \left[f_{2*}^{'}\right]_1$ denoting

$[f_2]_1\,(Q_1^*,\quad h_{1*}^{'}$ denoting $h_*^{'}$ denoting $h_1^{'}\,(Q_1^*)$, and $\left[g_*^{'}\right]_1$

Denoting $\left[g^{'}\right]_1(P_1^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_1]_1 + [m_a]_1 = -(\mu_1(a_1) + [f_{1*}]_1 + [f_{2*}]_1)m_1 - [f'_{1*}]_1[p_1^*](a_1)M_1 - [f'_{2*}]_1 p_1^*(a_1)W_1$$

$$m_1(t,0) = [h_*]_1 \int_0^\infty \beta_1(a_1)m_1(t,a_1)da_1 + \frac{[h_*^1]_1}{[h^*]_1}[p^*]_1(0)\,W_{,1} \qquad (30)$$

$$\frac{dW_1}{dt} = [g_*^1]_1[Q_1^*]M_1$$

Integrating $[m_1]_1(t,a_1)$ with respect to a from 0 to $\infty$, we have

$$\frac{dM_1(t)}{dt} = \int_0^\infty [mt]_1(t,a_1)da_1$$

$$= -\int_0^\infty [m_a]_1\,(t,a_1)da_1 - \int_0^\infty \mu(a_1)m(t,a_1)da_1$$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f'_{1*}]_1 \int_0^\infty p_1^*(a_1)da_1\right)M_1 - [f_{2*}^1]_1 \int_0^\infty p_1^*(a_1)da_1 W_1$$

$$= m_1(t,0) - \int_0^\infty \mu_1(a_1)\,m_1(t,a_1)da_1$$

$$-\left([f_{1*}]_1 + [f_{2*}]_1 + [f_{1*}']_1 P_1^*\right)M_1 - [f_{2*}']_1 P_1^* W_{,1}$$

From (30), we arrive at the equations

$$\frac{dW_1(t)}{dt} = \int_0^\infty ([h_*]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1$$

$$-\left([f_{1*}] + [f_{2*}]_1 + [f_{1*}']_1 P_1^*\right)M_1 + \left(\frac{[h_*']_1}{[h_*]^1}\,P_1^*(0) - f_{2*}' P_1^*\right)W_1,$$

(31)

$$\frac{dW_1(t)}{dt}\left[g^{'}{}_{*}\right]_1 Q_1^{*} M_{.1}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{*}]_1\beta_1(a_1) - \mu_1(a_1))m_1(t,a_1)da_1 = ([h_{*}]_1\beta_1(\theta_1) - \mu_1(\theta_1))M_{,1}$$

Where $0< \theta_1 = \theta_1(t) < \infty$.

And if $a_1^2\,[a_1^2]_1<4\,[a_2]_1<4\,a_2 e$,

Re $\lambda_1\left[\lambda_1\right]_1 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_1 = -\frac{1}{2}a_1[a2\ ]_1 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_{*}\,[h_{*}]_1\beta_1(a_1) - \mu_1(a_1)\right|$. Then$|n_1(t)|$ ⊠Δ∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium($\rho^{*}(a), \boldsymbol{Q}^{*}$) is asymptotically stable if either $a_1^2 <4\,a_2 e$ and $[a_1] > 2\Delta$,or $[a_1^2]$ ⊠4$[a_2]$e and $(a_1 - \sqrt{(a_1^2)_1 - 4(a_2)_1 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+1}\right)$ and $\beta_1(a_1) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_1(a_1) e^{-\int_0^a \mu_1(\tau_1) + d\tau_1 da} = \frac{1}{(m+\mu_0+1)^2} + \frac{2}{(m+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_1 + [\mu_0]$ 0.5. Assume$[f_1](x_1) = K_1 x_1$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_1(x_1) = [0x_1]/(1+x_1)$; the functional responses $[[f_2](x) = [0]\, x_1/(1+x_1)$and $h_1(x_1)$=(1+$\omega_1 x_1$)$/(1+x_1)$(See Holling [7] and Rashevsky [13])

Equation (14),now has the form of

$$\frac{m_1 + [\mu_0] + [K][P_0])}{[m_1]_1 + (\mu_0)_1 + [K][P_0])_1^3\ [K]_1 [P_0]_1^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1+P_0) < b_1$i e $[P_0] < b/(\delta_1 - b_1)$then the equilibrium$(P_0)$ $[P_0]_1(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_1$

$[f_{1*}] = K_1\ P^*\ [f_{2*}] = \frac{\sigma Q_1^*}{1+Q_1^*} \quad [h_*] = \frac{1+\omega_1 Q_1^*}{1+Q_1^*}$

$[f_{1*}] = K_1\ [f_{2*}^1] = \frac{\sigma_1}{(1+Q_1^*)^2}[\ g_*^1] = \frac{\sigma_1}{(1+P_1^*)^2} \quad [h_*^1] = \frac{\omega_1 - 1}{(1+Q_1^*)^2}$

## AN ANALOGY BETWEEN PREDATOR PREY AND ACCENTUATION ATTRION MODELS

**Best example how the model could be applied is the Schrödinger's equation itself wherein RHS gobbles up LHS, and the terms on RHS do the same with each other. It is like adding milk to water. Whether it happens instantaneously or time lag is there is a different question altogether dealt at the opportune points**

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| **Best example that can be given is the formulation of the quantum wave equation albeit in a nascent form for the recently found gravitational waves. This has been studied in length in various papers of gravitational waves, wherein the attribution of how the model is used is described. Note the calculation holds for only one module. This is in consideration to the fact that the relation is beteen the both debit systems and credit systems which as we have enucleated in the foregoing bears ample testimony for the application of the model** |
| **Gravitational quantum wave equation:** |
| **Equation Time-dependent equation (Wikipedia)** |
| **(1) In quantum mechanics, the Schrödinger equation is (=) a partial differential equation that describes how the quantum state of a quantum system changes with (e&eb) time. It was formulated in late 1925, and published in 1926, by the Austrian physicist Erwin Schrödinger.[1]** |
| **(2) In classical mechanics Newton's second law, (F = ma), is used to make (eb) a mathematical prediction as to what path a given system will take following a set of known initial conditions.** |
| **(3) In quantum mechanics, the analogue of Newton's law is (=) Schrödinger's equation for (e) a quantum system (usually atoms, molecules, and subatomic particles whether free, bound, or localised).** |
| **(4) It is not a simple algebraic equation, but in general a linear partial differential equation, describing (eb) the time-evolution of the system's wave function (also called a "state function").[2]:1–2** |
| **(5) The concept of a wavefunction is (=) a fundamental postulate of quantum mechanics.** |
| **(6) Although Schrödinger's equation is often presented as a separate postulate, some authors [3]: Chapter 3 show that some properties resulting from Schrödinger's equation may be deduced just from (e) symmetry principles alone, for example the commutation relations.** |
| **(7) Generally, "derivations" of the Schrödinger equation demonstrate (eb) its mathematical plausibility for describing (eb) wave–particle duality, but to date there are no universally accepted derivations of (e) Schrödinger's equation from (e) appropriate axioms.** |
| **(8) In the Copenhagen interpretation of quantum mechanics, the wave function is (=) the most complete description that can be given of (e) a physical system.** |
| **(9) Solutions to Schrödinger's equation describe (eb) not only molecular, atomic, and subatomic systems, but also macroscopic systems, possibly even the whole universe.[4]:292ff** |

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| **(10) The Schrödinger equation is consistent with (e&eb) classical mechanics but not with (e&eb) special relativity.** |
| **(11) Making quantum mechanics consistent with (e&eb) special relativity is one of the great achievements of (e) Quantum Field Theory and requires (e) the introduction of the creation and annihilation operators so that (e) the number of particles can no longer be considered constant.** |
| **(12) The Schrödinger equation is not the only way to make predictions in (eb) quantum mechanics—other formulations can be used, such as Werner Heisenberg's matrix mechanics, and Richard Feynman's path integral formulation** |
| **(13) The form of the Schrödinger equation depends on (=) the physical situation (see below for special cases).** |
| **(14) The most general form is the time-dependent Schrödinger equation, which gives a description of a system evolving with time: [5]:143** |
| **(15) A wave function that satisfies the non-relativistic Schrödinger equation with V = 0. In other words, this corresponds to (e &eb) a particle traveling freely through (e&eb) empty space. Time-dependent** |
| **Schrödinger equation (general)** |
| **(16) i ℏ ∂ ∂ t Ψ ( r , t ) = H ^ Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)={\hat {H}}\Psi (\mathbf {r} ,t) where i is the imaginary unit, ħ is the Planck constant divided by 2π, the symbol ∂/∂t indicates a partial derivative with respect to time t, Ψ (the Greek letter psi) is the wave function of (e) the quantum system, and Ĥ is the Hamiltonian operator (which characterizes the total energy of any given wave function and takes different forms depending on the situation).** |
| **(17) Each of these three rows is a wave function which satisfies (eb) the time-dependent Schrödinger equation for a harmonic oscillator. Left: The real part (blue) and imaginary part (red) of the wave function. Right:** |
| **(18) The probability distribution of finding the particle with (e&eb) this wave function at a given position.** |
| **(19) The top two rows are examples of stationary states, which correspond to (e&eb) standing waves.** |
| **(20) The bottom row is an example of a state which is (=) not a stationary state.** |
| **(21) The right column illustrates (eb) why stationary states are (=) called "stationary".** |

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| **(22) The most famous example is the non-relativistic Schrödinger equation for (e) a single particle moving in an electric field (but not a magnetic field; see the Pauli equation): [6]** |
| **Time-dependent Schrödinger equation** |
| **(Single non-relativistic particle)** |
| **(23) i ℏ ∂ ∂ t Ψ ( r , t ) = [ − ℏ 2 2 μ ∇ 2 + V ( r , t ) ] Ψ ( r , t ) {\displaystyle i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)} i\hbar {\frac {\partial }{\partial t}}\Psi (\mathbf {r} ,t)=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} ,t)\right]\Psi (\mathbf {r} ,t)** |
| **(24) Here μ is (=) the particle's "reduced mass", V is its potential energy, ∇2 is the Laplacian (a differential operator), and Ψ is (=) the wave function (more precisely, in this context, it is called the (=) "position-space wave function").** |
| **(25) In plain language, it means "total energy equals (=) kinetic energy plus potential energy", but the terms take unfamiliar forms for (e) reasons explained below.** |
| **(26) Given the particular differential operators involved, this is (=) a linear partial differential equation.** |
| **(27) It is also a diffusion equation, but unlike the heat equation, this one is (=) also a wave equation given (e&eb) the imaginary unit present in the transient term.** |
| **(28) The term "Schrödinger equation" can refer to (e) both the general equation (first box above), or the specific nonrelativistic version (second box above and variations thereof).** |
| **(29) The general equation is indeed quite general, used throughout quantum mechanics, for everything from the Dirac equation to quantum field theory, by plugging in various complicated expressions for the Hamiltonian. The specific nonrelativistic version is a simplified approximation to reality, which is quite accurate in many situations, but very inaccurate in others (see relativistic quantum mechanics and relativistic quantum field theory)** |
| **(30) To apply the Schrödinger equation, the Hamiltonian operator is set up for the system, accounting for the kinetic and potential energy of the particles constituting the system, then inserted into the Schrödinger equation. The resulting partial differential equation is solved for the wave function, which contains information about the system.** |
| **Time-independent equation** |

**(31) The time-dependent Schrödinger equation described above predicts that wave functions can form standing waves, called stationary states (also called "orbitals", as in atomic orbitals or molecular orbitals). These states are important in their own right, and if the stationary states are classified and understood, then it becomes easier to solve the time-dependent Schrödinger equation for any state. Stationary states can also be described by a simpler form of the Schrödinger equation, the time-independent Schrödinger equation. (This is only used when the Hamiltonian itself is not dependent on time explicitly. However, even in this case the total wave function still has a time dependency.)**

**Time-independent Schrödinger equation (general)**

**(32) E Ψ = H ^ ⊠ Ψ {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi } {\displaystyle E\Psi =\operatorname {\hat {H}} \Psi }**

**In words, the equation states:**

**(33) When the Hamiltonian operator acts on a certain wave function Ψ, and the result is proportional to the same wave function Ψ, then Ψ is a stationary state, and the proportionality constant, E, is the energy of the state Ψ.**

**(34) The time-independent Schrödinger equation is discussed further below. In linear algebra terminology, this equation is an eigenvalue equation.**

**(35) As before, the most famous manifestation is the non-relativistic Schrödinger equation for a single particle moving in an electric field (but not a magnetic field):**

**Time-independent Schrödinger equation (single non-relativistic particle)**

**(36) E Ψ ( r ) = [ − ℏ 2 2 μ ∇ 2 + V ( r ) ] Ψ ( r ) {\displaystyle E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )} E\Psi (\mathbf {r} )=\left[{\frac {-\hbar ^{2}}{2\mu }}\nabla ^{2}+V(\mathbf {r} )\right]\Psi (\mathbf {r} )**

**Gravitational waves: (Wikipedia)**

**(37) In theory, the loss of energy through gravitational radiation could eventually drop the Earth into the Sun. However, the total energy of the Earth orbiting the Sun (kinetic energy + gravitational potential energy) is about 1.14×1036 joules of which only 200 Watts (joules per second) is lost through gravitational radiation, leading to a decay in the orbit by about 1×10−15 meters per day or roughly the diameter of a proton. At this rate, it would take the Earth approximately 1×1013 times more than the current age of the Universe to spiral onto the Sun. This estimate overlooks the decrease in r over time, but the majority of the time the bodies are far apart and only radiating slowly, so the difference is unimportant in this example.**

**More generally, the rate of orbital decay can be approximated by[39]**

**(38) d r d t = − 64 5 G 3 c 5 ( m 1 m 2 ) ( m 1 + m 2 ) r 3 , {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,} {\displaystyle {\frac {\mathrm {d} r}{\mathrm {d} t}}=-{\frac {64}{5}}\,{\frac {G^{3}}{c^{5}}}\,{\frac {(m_{1}m_{2})(m_{1}+m_{2})}{r^{3}}}\ ,}**

**where r is the separation between the bodies, t time, G Newton's constant, c the speed of light, and m1 and m2 the masses of the bodies. This leads to an expected time to merger of [39]**

**t = 5 256 c 5 G 3 r 4 ( m 1 m 2 ) ( m 1 + m 2 ) . {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.} {\displaystyle t={\frac {5}{256}}\,{\frac {c^{5}}{G^{3}}}\,{\frac {r^{4}}{(m_{1}m_{2})(m_{1}+m_{2})}}.}**

**(39) For example, a pair of solar mass neutron stars in a circular orbit at a separation of 1.89×108 m (189,000 km) has an orbital period of 1,000 seconds, and an expected lifetime of 1.30×1013 seconds or about 414,000 years. Such a system could be observed by LISA if it were not too far away. A far greater number of white dwarf binaries exist with orbital periods in this range. White dwarf binaries have masses in the order of the Sun, and diameters in the order of the Earth. They cannot get much closer together than 10,000 km before they will merge and explode in a supernova which would also end the emission of gravitational waves. Until then, their gravitational radiation would be comparable to that of a neutron star binary.**

**(40) By substituting value of dr/dt in Schrödinger's wave equation we get the quantum wave equation for gravitational waves.**

In order to have (H3) and (H4) hold, $m, \mu_0$, K, b, δ, and $\omega$ need to satisfy

$$\frac{\omega_1}{(m+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1+b_1}+\sigma)^2}+\frac{2\omega_1}{([m_1]_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1+b_1}+\sigma_1)^3}<1$$

And

$$\frac{1}{(m_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1-b_1})^2}+\frac{2\omega_1}{(m_1+[\mu_0]_1+1+\frac{K_1 b_1}{\delta_1-b_1})^3}>1$$

When the inequalities (33) and (34) hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$(m_1 + [\mu_0]_1 + 1 + \frac{K_1 b_1}{\delta_1 - b_1} + (m_1 + [\mu_0]_1 + \frac{K_1 b_1}{\delta_1 - b_1} + \sigma_1)\, Q_1^*)^3$

$= m_1 + [\mu_0]_1 + 3 + \frac{K_1 b_1}{\delta_1 - b_1}(1 + Q^*)(1 + \omega_1 Q_1^*$

Now

$$\left[a_1^2\right] - 4\left[a_2\right] e = \left(\frac{2K_1 b_1}{\delta_1 - b_1}\right)^2 + \frac{4K_1 b_1 \sigma_1 Q_1^*}{(\delta_1 - b_1)(1 + Q_1^*)} + \frac{\sigma_1^2 Q_1^{*2}}{(1 + Q_1^*)^2}$$

$$\frac{4\, b_1 \sigma_1 (\delta_1 - b) Q^*}{\delta_1 (1 + Q_1^*)^2} - \frac{4\, (\delta_1 - b_1)^2 (1 - \omega_1) Q_1^*}{\delta_1 (1 + \omega_1 Q_1^*)(1 + Q_1^*)} P^*(0)$$

And

$\Delta = \lim_{0 < a < \infty} |[h_*]_1\, a e^{mu}_1 \quad [\mu_0]_1 - \frac{a_1}{a_1 + 1}$

Hence if $[a_1^2 > 4(a_2)_1 e$ and

$$[a_1] = \frac{2K_1 b_1}{\delta_1 - b_1} + \frac{\sigma_1 Q_1^*}{1 + Q_1^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_1]_1 = [f_{1*}]_1 + [f_{2*}]_1 + [f_{1*}] P_1^*$$

$$[a_2]_1 = \left[f^{1}_{2*}\right] P^{*}_{1} - \frac{h^{1}_{*}}{h_{*}} P^{*}(0)$$

$$e=[g^{1}_{*}]_1 Q^{*}_{1}$$

$$\eta_1(t) = \left[h_{*}\right]_1 \beta_1 \left(\theta_1(t)\right) - \mu_1(\theta(t)),$$

The system(31)can be written as the following system:

$$[M(t)]\left[a_1\right]M_1 - [a_2]_1 W_1 - [n(t)]M_1$$

$$w_1(t) = eM_1$$

Write

$$A=\begin{pmatrix} -[a_1]_1 - [a_2]_1 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_1 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_1 \\ W_1 \end{pmatrix}$$

The solution of (32)is

$$X_1(t) = e^{A,t}[X]_1(0) + \int_0 e^{At}$$

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# Presciential Dynamics And Contential Staticity Of Accentuation Attrition Models

## (Module Three)

**Abstract: Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-435) 1990. We demonstrate its applicability cosmology/ quantum computation economics et al. As stated in the foregoing we shall apply model to various cosmological problem and show the relevance to the same. . To exemplify model can be applied to the following situational significations: Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an**

**outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 Masaru Shibata and Takashi Nakamura** perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. **Authors** use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society **Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995**

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(g_3)Q_3$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_3(0) = 0, \quad g_3(P_3) > 0, \quad P_3 \text{ 0}, \qquad \lim_{P_3 \to \infty} g_3(P_3) = (K_3)_3 \text{ ∞}. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_3 + (\rho_a)_3 + D_3(a_3, P_3, Q_3)\rho_3 + 0, \qquad 0 < a_3 \text{ ∞},$$

$$\rho_3(t,0) = \int_0^\infty n_3(a_3, P_3, Q_3)\rho_3(t_3, a_3)da_3$$

Where $D_3$ is the mortality function and $n_3$ is the fertility function.

Put $_3$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_3}{dt}\right]_3 = x_3\left[\frac{dQ_2}{dt}\right]_3 - x_2\left[\frac{dQ_3}{dt}\right]_3$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_3 - x_4\left[\frac{dQ_2}{dt}\right]_3$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_3 - x_6\left[\frac{dQ_3}{dt}\right]_3$$

Where $x_3, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_3$ and $g_3(P_3)Q_3$

Now when you add the term $(b_3Q_3)$ with $g_3(P_3)(Q_3)$ we are infect see is that $(b_3Q_3)$ is credited incrementally part of $g_3(P_3)(Q_3)$ or instant nearly, when subtracted from $\frac{dQ_3}{dt}$ , it means that $g_3(P_3)(Q_3)$ is slowly eaten by $\frac{dQ_3}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_3 = \mu_3(a_3) + (f_3)_3(P_3)_3 + (f_2)_3(Q)_3 \quad (2)$$

$$n_3 = \beta_3(a_3)h_3(Q_3), \quad (3)$$

Where $\mu_3$the age- dependent debit rate whit is $\mu_3(a_3)$ 0 $\forall a > 0$; $(f_3)_3$ is the density- dependent death rate with $f_3(0) = 0$, $(f_3^3)_3(P_3)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_3)_3(P_3) = (k_3)_3$ ; $f_2$ is the functional response of

creditors to prey with $[f_2]_3(0) = 0, (f_2^3)_3(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_3\ (Q_3) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 3, h^3(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_3 + (\rho_a)_3 = -\mu_3(a_3) + (f_3)_3 P_3) + \big(f_2)_3(Q_3)\big)(\rho_3)_3 \quad (4)$$

$$\rho_3(t_3, 0) = h_3(Q_3)\int_0^\infty \beta_3\,(a_3)\rho_3(t, a_3)da, \quad (5)$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + g_3(P_3)Q_3 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_3$ and $Q_3$ of (4-6) satisfy

$$0 \ \rho_3(t, a_3) \ \rho_3^*(t, a_3),$$

$$0 \ Q_3(t) \ Q_3^*(t),$$

Where $\rho_3^*(t, a_3)$ and $Q_3^*(t)$, are solutions of

$$(\rho_t^*)_3 + (\rho_a^*)_3 = -(\mu_3(a_3) + (f_3)_3(P_3^*) + (f_2)_3 Q_3^*))\,\rho_3^*, \quad (7)$$

$$\rho_3^*\,(t, 0) = \int_0^\infty \beta_3(a_3)\,\rho_3^*\,(t, a_3)da_3 \quad (8)$$

$$\frac{dQ_3^*}{dt} = -\,b_3\,Q_3^* + g_3(P_3^*)Q_3^*, \quad (9)$$

With $P_3^*(t) = \int_0^\infty \rho_3^*\ (t, a_3)da_3$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_3^*(t)}{dt} = P_3^*(t)\big(r_3 - (f_3)(P_3^*) - (f_2)_3(Q^*)\big), \quad (10)$$

$$\frac{dQ_3^*(t)}{dt} = Q_3^*(t)(-b_3 + g_3(P_3^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_3(a)e^{-\int_0^a \mu_3(\tau)dt}e^{-ra}\ da_3 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_3^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(t)dt}\ da_3 > 1, \quad (H1)$$

Which guarantees $r_3 > 0$.

*Remark.* If $\beta_3(a) = [\beta_0]_3 a_3 e^{-\alpha_3 a_3}, \mu_3(a_3) \equiv \mu_3$ then the solution of (12) is

$$r_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_3]_3(P_3) = 0, [f_2](Q_3) = \lambda_3 Q_3\ h_3(Q_3) = 1, \quad and\ \ g_3(P_3) = c_3 P_3$$

$$[\rho_3]_3 + [\rho_a]_3 = -(\mu_3 + \lambda_3 Q_3)\rho_3$$

$$\rho_3(t,0) = \int_0^{\infty} [\beta_0]_3 a_3 e^{-\alpha_3 a_3} \rho_3(t,a_3)\, da_3$$

$$\frac{dQ_3}{dt} = -b_3 Q_3 + c_3 P_3 Q_3$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_3}{dt} = \omega_3 P_3 - \lambda_3 P_3 Q_3$$

$$\frac{dQ_3}{dt} = b_3 Q_3 - c_3 P_3 Q_3$$

Where $\omega_3 = \sqrt{\beta_3} - \alpha_3 - \mu_3$

## THE EQUILIBRIUM $(\rho_0)_3(a_3), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_3(a_3), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_3 + (\rho_a)_3 = -\big(\mu_3(a_3) + (f_3)_3(P_3)\big)\rho,$$

$$\rho_3(t,0)$$

$$= \int_0^\infty \beta_3\,(a_3)\rho_3(t,a_3)da_3 \qquad (13)$$

The solution of (13) has the form of

$$\rho_3(t,a_3) = Ce_3^{r(t-a_3)} - \int_0^a (\tau)d\tau - \int_0^3 (f_3)_3\big(P_3(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_3(t) = C_3 e^{r_3 t - \int_0^t (f_3)_3(P_3(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_3^{-3}(r)$ if $k_3 > r$. This is equivalent to

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{(k_3)_3 a_3} da_3 < 1.$$

In fact, if we define

$$R_3(s_3) = \int_0^\infty \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau)d\tau} e^{-s_3 a_3}\, da_3$$

$R(s)$ is a decreasing function. $R_3(r_3) = 1$ And $R_3(k_3)_3 < 1$ imply $r_3 < (k_3)_3$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_3(a_3), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{k_3 a_3} da_3$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_3 \equiv \int_0^\infty (\rho_0)_{33}(a_3)da_3$ *satisfies*

$$\int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau)d\tau} e^{-(f_3)_3(P_0)a_3} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_3(a_3), 0)$

Let $u_3(t, a_3) = \rho(t, a_3) - \rho_0(a_3)$ Then

$$N_3(t) \equiv \int_0^{\infty} u_3(t, a_3) da_3 = \int_0^{\infty} \rho_3(t, a_3) da_3 - \int_0^{\infty} (\rho_0)_3(a_3) da = P_3(t) - P_0)_3$$

And

$$(u_3)_3 + (u_a)_3 = (\rho_t)_3 + (\rho_a)_3 - (\rho_0^3)_3(a_3)$$

$$= -\mu_3(a_3) + (f_3)_3(P_3) + (f_2)_3(Q_3))(u_3 + (\rho_0)_3) + \mu_3(a_3)(\rho_0)_3 + (f_3)_3(P_0)_3(\rho_0)_3$$

$$= -\mu_3(a_3)u_3 - (f_3)_3(P_3) + (f_2)_3(Q_3))u_3 - (\rho_0)_3 (f_3)_3(P_3) - (f_3)_3(P_0)_3$$

$$+ (f_2)_3(Q_3)), \qquad (15)$$

Where $(\rho_0^3)_3(a_3)$ denotes $d(\rho_0)_3(a_3)/da_3$

Expansion $(f_3)_3(P_3)$ and $(f_2)_3(Q_3)$ at $(P_0)_3$ and 0 yield respectively

$$(f_3)_3(P_3) = (f_3)_3((P_0)_3 + N_3)$$

$$= (f_3)_3(P_0)_3 + (f_3^3)_3(P_0)_3 N_3 + o_3(N_3), \qquad (16)$$

$$(f_2)_3(Q_3)$$

$$= (f_2^3)_3(0)Q + o_3(Q_3) \qquad (17)$$

Write $f_3(P_0) = f_{30}, f_3'(P_0) = f_{30}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_3)_3 + (u_a)_3 \approx -(\mu_3(a_3) + (f_{30})_3)u - (f_{30}^3)_3(\rho_0)_3(N_3) - (f_{20}^3)_3(\rho_0)_3(Q_3)$$

$$u_3(t, 0) = \rho_3(t, 0)(\rho_0)_3(0)$$

$$= h_3(Q_3)\int_0^{\infty} \beta_3(a_3)(u_3(t,a) + (\rho_0)_3(a_3)da_3 - (\rho_0)_3\,(0)$$

$$= h_3(Q_3)\int_0^{\infty} \beta_3(a_3)(u_3(t,a_3)da_3 + (h_3(Q_3) - 3)(\rho_0)_3\,(0) \qquad (18)$$

Expanding

$$h_3(Q_3) = 1 + \ (h^3)_3(0)Q_3 + 0(Q_3) \equiv 1 + h_0^{'}Q_3 + 0Q_3$$

And substituting it into (38), we arrive at

$$u_3(t,0) \approx \int_0^{\infty} \beta_3\ (a_3)u_3(t,a_3)da_3 + (h_{a_3}^{'})(\rho_0)_3Q_3$$

The linearization is the following :

$$[u_3] + [u_3] = -(\mu_3(a_3) + [t_{30}]_3u_3 - [\alpha_3]_3(a_3)N_3 - [\alpha_2]_3(a_3)Q_3$$

$$u_3(t,0) = \int_0^{\infty} \beta_3(a_3)u_3(t,a_3)da_3 + [\xi_3]Q_3 \qquad (19)$$

$$\frac{dQ_3}{dt} = -b_3Q_3$$

where

$$[\alpha_3]_3(a_3) = [t_3^3a]_3[p,o]_3(a_3)d2(a_3) = [f_{20}^3]_3(\rho\ o\ )_3(a_3)\xi = [\mathrm{h\ o}^3]_3[\rho\mathrm{o}]_3(\mathrm{o}) \text{ and } \mathrm{b}_3 = \mathrm{b}_3 - g_3(P_3)$$

If $g_3(P_0) > b$, clearly, $(\rho_0)_3(a_3), 0)$ is unstable and the predator population is persistent. Assume $g_3(P_0)_3 < b$.

Then, $Q_3(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_3 + u_a = -(\mu_3(a_3) + (f_{30})_3)u_3 - (\alpha_3)_3(a_3)N_3$$

$$= -(\mu_3(a_3) + f_{30})u_3 - \alpha_3(a_3)\int_0^\infty u_3(t, a_3)da_3$$

$$u_3(t, 0) = \int_0^\infty \beta_3\,(a_3)u_3(t, a_3)da_3$$

Using the method developed by Cushing [2, 3], we suppose

$$u_3(t, a_3) = \tilde{u}_3(a_3)^{e^{\gamma(t-a_3)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_3}\,da_3$. Then,

A solution of the first differential equation for $\tilde{u}(a_3)$ gives,

$$\tilde{u}_3(a_3) = e^{-\int_0^a \mu_3(\tau_3)d\tau - (f_{30})_3 a}\left(\tilde{u}_3(0) - w\int_0^a e^{\int_0^\tau \mu_3(s_3)ds + (f_{30})}\tau_3\,(\alpha_3)_3(\tau_3)e^{\gamma_3\tau}d\tau\right).$$

Substituting the expression (13) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\right)\tilde{u}_3(0)$$

$$+w\int_0^\infty \beta_3(a_3)e^{-\gamma_3 a_3 - (f_{30})_3 a_3 - \int_0^a \mu_3(\tau_3)d\tau}\int_0^a (\alpha_3)_3(\tau_3)e^{-\gamma_3\tau_3}e^{(f_{30})_3\tau_3 + \int_0^\tau \mu_3(s_3)ds}d\tau_3\,da_3$$

$$= 0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(\tau_3)d\tau} da\right)(\tilde{u}_3(0))_3$$

$$+\left(-3-\int_0^\infty e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(\tau_3)d\tau}\int_0^a e^{(f_{30})_3\tau_3+\int_0^\tau \mu_3(s_3)ds+y\tau}\alpha_3(\tau)d\tau\, da\right)w=0. \quad (23)$$

Define $Y(\gamma,a)=e^{-\gamma_3 a_3-(f_{30})_3 a_3-\int_0^a \mu_3(t_3)d\tau}$. Then

$$\left(1-\int_0^\infty \beta_3\,(a_3)Y_3(\gamma_3,a_3)da_3\right)u_3(0)$$

$$+w\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)\int_0^a(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\, da_3=0, \quad (24)$$

$$\left(\int_0^\infty Y_3(\gamma_3,a_3)da_3\right)\tilde{u}_3(0)$$

$$+\left(-1-\int_0^\infty Y_3(\gamma_3,a_3)\int_0^a(\alpha_3)_3(\tau)_3(Y^{-3})_3(\gamma_3,\tau_3)d\tau_3\, da_3\right)w=0.$$

Equqtion (20) has a solution of the form of $u_3(t,a_3)=u_3(a_3)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_3(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_3(\gamma_3,a_3)=\int_0^a(\rho_0)_3(\tau_3)Y^{-3}\,(\gamma_3,a_3)d\tau_3$. We obtain the characteristic equation

$$f_{30_3}^l\left(\int_0^\infty Y_3(\gamma_3,a_3)da_3\right)\left(\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$+\left(1-\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)da\right)\left(1+(f_{30}^l)_3\int_0^\infty Y_3(\gamma_3,a_3)I_3(\gamma_3,a_3)da_3\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_3(a_3)=(\rho_0)_3(0)e^{-\int_0^a \mu_3(\tau_3)d\tau_3-(f_{30})_3a_3}$$

$$I_3(\gamma_3,a_3)=\int_0^a (\rho_0)_3(\tau_3)e^{\gamma_3\tau_3+(f_{30})_3\tau+\int_0^\tau \mu_3(s_3)ds_3}d\tau$$

$$=\frac{(\rho_0)_3(0)}{\gamma_3}(e^{\gamma_3a_3}-1)$$

Let

$$E_3(a_3)=e^{-(f_{30})_3a_3-\int_0^a \mu_3(\tau_3)d\tau_3}$$

Then $[Y_3]_3[\gamma_3]_3,a_3)=E_3(a_3)e^{-\gamma_3a_3}$. by a simple calculation,

$$\int_0^\infty \beta_3(a_3)Y_3(\gamma_3,a_3)I(\gamma_3,a_3)da=\frac{[\rho_0]_3(0)}{\gamma}\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3a_3}\ da\right);$$

$$\int_0^\infty Y_3(\gamma_3,a_3)I(\gamma_3,a_3)da=\frac{[\rho_0]_3(0)}{\gamma_3}\left(\int_0^\infty E_3(a_3)e^{-\gamma_3a_3}\ da_3\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_3(0)}{\gamma_3}[f_{30}^3]\int_0^\infty E_3(a_3)e^{-\gamma_3a_3}da_3\ (\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3a_3}da_3)$$

$$+\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right).$$

$$\left(1+\frac{(\rho_0)_3(0)}{\gamma_3}(f_{30}^{'})_3\left(\int_0^\infty E_3(a_3)da_3-\int_0^\infty E_3(a_3)e^{-\gamma_3 a_3}da_3\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3\right)(1+\frac{(f_{30}^{'})_3}{\gamma_3}(\rho_0)_3(0)\int_0^\infty E_3(a_3)da_3)=0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_3(a_3)Y_3(\gamma_3 a_3)da_3 = 1-\int_0^\infty \beta_3(a_3)E_3(a_3)e^{-\gamma_3 a_3}da_3 \neq 0,$$

Since $\int_0^\infty \beta_3(a_3)Y_3(\gamma_3 a_3)\int_0^a \alpha_3(\tau_3)Y^{-3}(\gamma_3,\tau_3)d\tau_3 da_3 > 0$. Therefore, the characteristic solution is

$$\gamma_3 = -(f_{30}^{'})_3(\rho_0)_3\int_0^\infty E_3(a_3)da_3.$$

Since $f_3^{'}(P_0)_3 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_3), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_3), 0)$ which is related to the extinction of debit summation types

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_3 + (\mu_3(a_3) + [f_3]_3(P_3^*) + [f_2]_3(Q_3^*))P_3^* = 0, \quad (26)$$

$$p_3^*(0) = h_3(Q_3^*)\int_0^\infty \beta_3(a_3)p_3^*(a_J)p_3^*(a_3)da_3 \quad (27)$$

$$g_3(P_3^*) - b_3 = 0, \quad (28)$$

Where $P_3^* = \int_0^\infty p_3^*(a_3)\, da_3$

Equation (28) yields $P^* = g^{-3}(b)$,

From (26),

$$p_3^*(a) = p_3^*(0)e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}$$

Substituting into (27), we need to solve

$$\left(1 - h_3(Q_3^*)\int_0^\infty \beta_3(a_3)e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(Q_3^*)a_3}da_3\right)p_3^*(0) = 0$$

For $Q_3^*$

Fix $p_3^*(0)$ and $P_3^*$ Let

$$H_3(z_3) = 1 - h_3(z_3)\int_0^{\infty} \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3} da_{,3}$$

And assume that

$$[k_3]_3 \int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(g_3^{-3}(b_3))a_3 - [k_2]_3 a_3}\, da_3 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_3(z_3)] > 0.$$

Clearly,

$$\frac{dH_3}{dz_3} = h_3(z)([f_2^{'}]_3(z_3))\int_0^{\infty} a_3\beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - f_3(P_3^*)a_3 - f_2(z_3)a_3} da_{3}$$

$-h^{'}(z_3)\int_0^{\infty} \beta_3(a_3)\, e^{\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3 - [f_2]_3(z_3)a_3}\ da_3$

By the assumptions on $\left[f_2^3\right]_3 and\ [h^3]^3 dH_3/dz > 0$. Hence, if

$H_3(0) = 1 - \int_0^{\infty} \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3(P_3^*)a_3} da_3 < 0,$ (29)

There exists a unique positive solution $Q^*$

Now, we claim that $H_3(0) < 0$ follows if $g_3^{-}(b_3) < [f_3^{-3}]_3(r_3)$, where $r_3$ is the solution of (12).

Indeed, since

$$R_3(s_3) = \int_0^{\infty} \beta_3(a_3)\, e^{-\int_0^a \mu_3(\tau_3)d\tau_3 e^{-y\frac{a}{3}}da_3}$$

Is a decreasing function of $s_3$ when $P_3^* = [g^{-3}]_3(b_3) < [f_3^{-3}]_3(r_3)$, that is $[f_3]_3(P_3^*) < r$, $R\big(f_3(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_3(g^{-3}(b))) > 1, i.e.,$

$\int_0^{\infty} \beta_3(a_3)e^{-\int_0^a \mu_3(\tau_3)d\tau_3 - [f_3]_3\left(g_3^{-3}(b)\right)a_3} > 1.$ (H4)

Then the system (4)-(6) has a unique positive equilibrium $(P_3^*(a_3), Q_3^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_3^*(a_3), Q_3^*)$ is more complicated. We again consider perturbations $m_3(t_3a_3) = p_3(t_3, a_3) - p_3^*(a_3)$ and $W_3(t) = Q_3(t) - Q^*$.Then

$M_3(t) \equiv \int_0^{\infty} m_3(t, a_3)da_3 = P_3(t) - P_3^*,$

$$[m_3]_3 + [m_a]_3 = [p_3]_3 + [p_a]_3 - \left[p^*\right]_3(a_3) \approx -(\mu_3(a_3) + [f_{3*}]_3 + [f_{2*}]_3)m_3$$

$$-[f_{3*}^3]_3 p_3^*(a_3)M - [f_{2*}^3]_3 p_3^*(\mathrm{a})W_3,$$

$$m_3(t,0)p_3(t,0) - p_3^*(0) \approx h_* \int_0^{\infty} \beta_3(a_3)m_3(t,a_3)\, da + \frac{[h_*^3]_3}{[h_*]_3}\, p_3^*(0)W_3$$

$$\frac{dW_3}{dt} = \frac{dQ_3}{dt} \approx [g_*^3]_3 Q_3^* M_{,3}$$

Where $[f_{3*}]_3$ denoting $[f_3(P^*)]_3$, $[f_{3*}]_3$ denoting $[f_2(P^*)]_3$, $[f'_{3*}]_3$ denoting $[f'_3(P^*)]_3 [f'_{2*}]_3$ denoting $[f_2]_3$ $(Q_3^*,\quad h_3^*$ denoting $h'_*$ denoting $h'_3$ $(Q_3^*)$, and $[g'_*]_3$

Denoting $[g']_3(P_3^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_3]_3 + [m_a]_3 = -(\mu_3(a_3) + [f_{3*}]_3 + [f_{2*}]_3)m_3 - [f'_{3*}]_3 [p_3^*](a_3)M_3 - [f'_{2*}]_3 p_3^*(a_3)W_3$$

$$m_3(t,0) = [h_*]_3 \int_0^\infty \beta_3(a_3)m_3(t,a_3)da_3 + \frac{[h_*^3]_3}{[h^*]_3}[p^*]_3(0)\,W_{,3} \qquad (30)$$

$$\frac{dW_3}{dt} = [g_*^3]_3[Q_3^*]M_3$$

Integrating $[m_3]_3(t,a_3)$ with respect to a from 0 to ∞, we have

$$\frac{dM_3(t)}{dt} = \int_0^\infty [mt]_3(t,a_3)da_3$$

$$= -\int_0^\infty [m_a]_3\,(t,a_3)da_3 - \int_0^\infty \mu(a_3)m(t,a_3)da_3$$

$$-\left([f_{3*}]_3 + [f_{2*}]_3 + [f'_{3*}]_3 \int_0^\infty p_3^*(a_3)da_3\right)M_3 - [f_{2*}^3]_3 \int_0^\infty p_3^*(a_3)da_3 W_3$$

$$= m_3(t,0) - \int_0^\infty \mu_3(a_3)\, m_3(t,a_3)da_3$$

$$-\left([f_{3*}]_3 + [f_{2*}]_3 + [f'_{3*}]_3 P_3^*\right)M_3 - [f'_{2*}]_3 P_3^* W_{,3}$$

From (30), we arrive at the equations

$$\frac{dW_3(t)}{dt} = \int_0^{\infty} ([h_*]_3\beta_3(a_3) - \mu_3(a_3))m_3(t,a_3)da_3$$

$$-\left([f_{3*}] + [f_{2*}]_3 + \left[f_{3*}^{'}\right]_3 P_3^*\right) M_3 + \left(\frac{[h_*^{'}]_3}{[h_*]^3}\, P_3^*(0) - f_{2*}^{'} P_3^*\right) W_3,$$

(33)

$$\frac{dW_3(t)}{dt}\left[g^{'}{}_*\right]_3 Q_3^*\, M_{.3}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{"*}]_3\beta_3(a_3) - \mu_3(a_3))m_3(t,a_3)da_3 = ([h_*]_3\beta_3(\theta_3) - \mu_3(\theta_3))M_{,3}$$

Where $0< \theta_3 = \theta_3(t) < \infty$.

And if $a_3^2\,[a_3^2]_3<4\,[a_2]_3<4\,a_2 e$,

$\mathrm{Re}\,\lambda_3\left[\lambda_3\right]_3 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_3 = -\frac{3}{2}a_3[a2\ ]_3 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_*\,[h_*]_3\beta_3(a_3) - \mu_3(a_3)\right|$. Then$|n_3(t)|$ ⊠⊿∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_3^2 <4\,a_2 e$ and $[a_3] > 2\Delta$,or $[a_3^2]$ ⊠$[a_2]e$ and $(a_3 - \sqrt{(a_3^2)_3 - 4(a_2)_3 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+3}\right)$ and $\beta_3(a_3) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^\infty \beta_3(a_3) e^{-\int_0^a \mu_3(\tau_3) + d\tau_3 da} = \frac{3}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_3 + [\mu_0]$ 0.5. Assume$[f_3](x_3)= K_3 x_3$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_3(x_3)= [0x_3]/(3 + x_3)$; the functional responses $[[f_2](x) = [0]\, x_3/(1 + x_3)$and $h_3(x_3)=(1+\omega_3 x_3)/(1 + x_3)$(See Holling [7] and Rashevsky [13])

Equation (34),now has the form of

$$\frac{m_3 + [\mu_0] + [K][P_0])}{[m_3]_3 + (\mu_0)_3 + [K][P_0])_3^3\ [K]_3 [P_0]_3^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(3 + P_0) < b_3$i e $[P_0] < b/(\delta_3 - b_3)$then the equilibrium$(P_0)$ $[P_0]_3(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_3$

$[f_{3*}] = K_3\ P^*\ [f_{2*}] = \frac{\sigma Q_3^*}{3+Q_3^*}$ $[h_*] = \frac{3+\omega_3 Q_3^*}{3+Q_3^*}$

$[f_{3*}] = K_3\ [f_{2*}^3] = \frac{\sigma_3}{(3+Q_3^*)^2}[\ g_*^3] = \frac{\sigma_3}{(3+P_3^*)^2}$ $[h_*^3] = \frac{\omega_3 - 3}{(3+Q_3^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0, K, b, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_3}{(m+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma)^2}+\frac{2\omega_3}{([m_3]_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3+b_3}+\sigma_3)^3}<1$$

And

$$\frac{1}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^2}+\frac{2\omega_3}{(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}+(m_3+[\mu_0]_3+\frac{K_3b_3}{\delta_3-b_3}+\sigma_3)\,Q_3^*)^3$$

$$=m_3+[\mu_0]_3+1+\frac{K_3b_3}{\delta_3-b_3}(1+Q^*)(1+\omega_3Q_3^*$$

Now

$$\left[a_3^2\right]-4\left[a_2\right]e=\left(\frac{2K_3b_3}{\delta_3-b_3}\right)^2+\frac{4K_3b_3\sigma_3\,Q_3^*}{(\delta_3-b_3)(1+Q_3^*)}+\frac{\sigma_3^2\,Q_3^{*2}}{(1+Q_3^*)^2}$$

$$\frac{4\,b_3\sigma_3\,(\delta_3-b)Q^*}{\delta_3\,(1+Q_3^*)^2}-\frac{4\,(\delta_3-b_3)^2(1-\omega_3)Q_3^*}{\delta_3(1+\omega_3Q_3^*)(1+Q_3^*)}P*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_3\,ae_3^{mu}\;\;[\mu_0]_3-\frac{a_3}{a_3+3}$$

Hence if $[a_3^2 > 4(a_2)_3 e$ and

$$[a_3] = \frac{2K_3 b_3}{\delta_3 - b_3} + \frac{\sigma_3 Q_3^*}{1 + Q_3^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_3]_3 = [f_{3*}]_3 + [f_{2*}]_3 + [f_{3*}] P_3^*$$

$$[a_2]_3 = \left[f_{2*}^3\right] P_3^* - \frac{h_*^3}{h_*} P^*(0)$$

$$e = [g_*^3]_3 Q_3^*$$

$$\eta_3(t) = \left[h_*\right]_3 \beta_3 \left(\theta_3(t)\right) - \mu_3(\theta(t)),$$

The system(33)can be written as the following system:

$$[M(t)]\left[a_3\right] M_3 - [a_2]_3 W_3 - [n(t)] M_3$$

$$w_3(t) = eM_3$$

Write

$$A = \begin{pmatrix} -[a_3]_3 - [a_2]_3 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_3 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_3 \\ W_3 \end{pmatrix}$$

The solution of (32)is

$$X_3(t) = e^{A,t}[X]_3(0) + \int_0 e^{At}$$

**This equation is useful in following type of problems . For actual model kindly see various papersThere is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant Λ, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive Λ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the Λ>0 generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible.**

**Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of (e) GW signals from BBHs with component spins $\chi_{1,2} \in [-0.99, 0.99]$ aligned (e&eb) with the orbital angular momentum, component masses $m_{1,2} \in [2,48] M_\odot$, and (e&eb) total mass $M_{total} \leq 50\ M_\odot$. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins >+0.6 and component masses $\gtrsim 20\ M_\odot$. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201 STOP Advanced LIGO may be the first experiment to (e) detect gravitational waves. Through superradiance of stellar black holes, it may also be (=) the first experiment to discover the QCD axion with (e&eb) decay constant above the GUT scale. When an axion's Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e) the black hole, forming a "gravitational atom." Through the superradiance process, the number of axions occupying (e) the bound levels grows exponentially, extracting (e) energy and angular momentum from the black hole. Axions**

**transitioning between levels of the gravitational atom and (e&eb) axions annihilating to (e) gravitons can produce (eb) observable gravitational wave signals. The signals are long-lasting, monochromatic, and can be distinguished from (e) ordinary astrophysical sources. Authors estimate up to O (1) transition events at aLIGO for (e) an axion between 10^-11 and (e&eb) 10^-10 eV and up to 10^4 annihilation events for (e) an axion between 10^-13 and 10^-11 eV. In the event of a null search, aLIGO can constrain (e) the axion mass for (e) a range of rapidly spinning black hole formation rates. Axion annihilations are also promising for (e) much lighter masses at future lower-frequency gravitational wave observatories; the rates have (e) large uncertainties, dominated by (e) supermassive black hole spin distributions. Projections for aLIGO are (=) robust against perturbations from (e) the black hole environment and account for (e) updated exclusion on (e&eb) the QCD axion of 6*10^-13 eV < ma < 2*10^-11 eV suggested by (e) stellar black hole spin measurements. Subjects: High Energy Physics - Phenomenology (hep-ph); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 91, 084011 (2015) DOI: 10.1103/PhysRevD.91.084011Cite as:arXiv:1411.2263 [hep-ph] (or arXiv:1411.2263v3 [hep-ph] for this version) Discovering the QCD Axion with Black Holes and Gravitational Waves Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang STOP Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric**

**charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by (e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave**

**signal. The signal sweeps upwards in frequency from 35 to 250 Hz with (e&eb) a peak gravitational-wave strain of (e) 1.0×10−21. It matches the waveform predicted by general relativity for (e) the inspiral and merger of a pair of black holes and the ringdown of the resulting (eb) single black hole. The signal was observed with a matched-filter signal-to-noise ratio of (e) 24 and a false alarm rate estimated to be (=) less than 1 event per 203 000 years, equivalent to (=) a significance greater than 5.1σ. The source lies at (eb) a luminosity distance of 410+160−180 Mpc corresponding to (e&eb) a redshift z=0.09+0.03−0.04. In the source frame, the initial black hole masses are (=) 36+5−4M⊙ and 29+4−4M⊙, and the final black hole mass is (=) 62+4−4M⊙, with (e&eb) 3.0+0.5−0.5M⊙c2 radiated in (eb) gravitational waves. All uncertainties define (eb) 90% credible intervals. These observations demonstrate (eb) the existence of binary stellar-mass black hole systems. Observation of Gravitational Waves from a Binary Black Hole Merger B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 061102 – Published 11 February 2016 This is the first direct detection of gravitational waves and (e&eb) the first observation of a binary black hole merger. Received 21 January 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.061102 STOP This is the second in a series of three papers on (e&eb) systematic analysis of rank 1 Coulomb branch geometries of (e) four dimensional N=2 SCFTs. In the first paper authors developed a strategy for (e) classifying physical rank-1 CB geometries of N=2 SCFTs. They show how to carry out this strategy computationally to construct (eb) the Seiberg-Witten curves and (e&eb) one-forms for all the rank-1 SCFTs. Explicit expressions are given for all cases, with the exception of the Nf=4 SU (2) gauge theory and (e&eb) the En SCFTs which were previously constructed. Classification includes all known rank-1 theories plus (+) a new one with (e&eb) an abelian flavor group, plus (+) nine additional theories whose existence is more speculative. Four of those, reported in first paper, depend on (=) the assumption of new frozen rank-1 SCFTs. Here authors also show that the assumption of the existence of certain rank-0 N=2 SCFTs leads to (eb) five additional consistent**

**rank-1 CB geometries. Subjects: High Energy Physics - Theory (hep-th) Cite as:arXiv: 1601.00011 [hep-th] (or arXiv: 1601.00011v2 [hep-th] for this version Geometric constraints on the space of N=2 SCFTs II: Construction of special Kähler geometries and RG flows Philip C. Argyres, Matteo Lotito, Yongchao Lü, Mario Martone (Submitted on 31 Dec 2015 (v1), last revised 6 Jan 2016 (this version, v2)) STOP Authors propose a model of Dark Energy in which the field currently dominating (e) the energy density of the universe is (=) an "axion field" linearly coupled to (e&eb) the Pontryagin density, tr (FΛF), (i.e., the exterior derivative of (e) the Chern-Simons form) of (e) a massive gauge field. Authors assume that the axion has (e) self-interactions corresponding to (e&eb) a non-trivial (exponential) potential. Authors argue that a non-vanishing magnetic helicity of (e) the gauge field triggers (eb) slow-rolling of the axion at field values far below (eb) the Planck scale. Proposal leads to (eb) a "Tracking Dark Energy Scenario" in which the contribution of the axion energy density to (e) the total energy density is (=) constant (and small) during (e&eb) the early radiation phase, until (e) a secular growth term proportional to (e) the Pontryagin density of (e) the gauge field becomes (=) dominant. The initially small contribution of (e) the axion field to (e) the total energy density is related to the observed small baryon-to-entropy ratio. Some speculations concerning the nature of the gauge field are offered.**

**Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Cite as:arXiv: 1601.00057 [hep-th] (or arXiv:1601.00057v2 [hep-th] for this version) Tracking Dark Energy from Axion-Gauge Field Couplings Stephon Alexander, Robert Brandenberger, Juerg Froehlich (Submitted on 1 Jan 2016 (v1), last revised 23 Jan 2016 (this version, v2)) STOP Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a**

**pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version) Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP In this work authors obtain bounds on (e&eb) the topological Abelian string-vortex and on (e&eb) the string-cigar, by using (e) a new measure of configurational complexity, known as (=) configurational entropy. In this way, the information-theoretical measure of (e) six-dimensional braneworld scenarios is capable to (e) probe situations where (e) the parameters responsible for (e) the brane thickness are arbitrary. The so-called configurational entropy (CE) selects (e&eb) the best value of the parameter in (eb) the model. This is accomplished by (e) minimizing the CE, namely, by (e) selecting the most appropriate parameters in the model that correspond to (e&eb) the most organized system, based upon (e) the Shannon information theory. This information-theoretical measure of complexity provides (eb) a complementary perspective to (e) situations where (e) strictly energy-based arguments are (=) inconclusive. Authors show that the higher the energy the higher the CE, what shows (eb) an important correlation between the energy of the localized field configuration and (e&eb) its**

**associated entropic measure. Subjects: High Energy Physics - Theory (hep-th) DOI:10.1016/j.physletb.2016.02.038 Cite as: arXiv: 1601.00076 [hep-th] (or arXiv: 1601.00076v2 [hep-th] for this version) Bounds on topological Abelian string-vortex and string-cigar from information-entropic measure R. A. C. Correa, D. M. Dantas, C. A. S. Almeida, Roldao da Rocha (Submitted on 1 Jan 2016 (v1), last revised 17 Feb 2016 (this version, v2)) STOP Sum of the Holevo quantity (that bounds the capacity of quantum channels to transmit (e&eb) classical information about an observable) and (e&eb) the quantum discord (a measure of the quantumness of (e) correlations of that observable) yields (eb) an observable-independent total given by (e) the quantum mutual information This split naturally delineates ,dilates and disseminates information about (e&eb) quantum systems accessible to (e) observers -- information that is redundantly transmitted by (e&eb) the environment -- while showing that it is (=) maximized for (e) the quasi-classical pointer observable. Other observables are accessible only via (e&eb) correlations with the pointer observable .Further authors prove an anti-symmetry property relating accessible information and (e&eb) discord. It shows that information becomes (=) objective -- accessible to (e) many observers -- only as quantum information is relegated to (e) correlations with (e&eb) the global environment, (eb) and, therefore, locally inaccessible. The resulting complementarity explains (eb) why, in a quantum Universe, we perceive objective classical reality while (e) flagrantly quantum superpositions are out of reach. Subjects: Quantum Physics (quant-ph) Journal reference: Scientific Reports 3, 1729 (2013) Report number: LAUR-13-26386 Cite as:arXiv: 1303.4659 [quant-ph] (or arXiv: 1303.4659v2 [quant-ph] for this version) Complementarity of quantum discord and classically accessible information Michael Zwolak, Wojciech H. Zurek STOP Niels Bohr proposed that the outcome of the measurement becomes (=) objective and real, and, hence, classical, when (e) its results can be communicated by (e) classical means. In this work authors revisit Bohr's postulate using (e) modern tools from (e) the quantum information theory. They find a**

**full confirmation of Bohr's idea: if a measurement device is in a nonclassical state, (eb) the measurement results cannot be communicated perfectly by classical means. In this case some part of information in the measurement apparatus is lost (e) in the process of communication: the amount of this lost information turns out to be (=) the quantum discord. The information loss occurs even when (e) the apparatus is not entangled with the system of interest. The tools presented in this work allow (eb) to generalize Bohr's postulate: authors show that for pure system-apparatus states (eb) quantum communication does not provide any advantage when (e) measurement results are communicated to (e) more than one recipient. They further demonstrate the superiority of quantum communication to (e&eb) two recipients on (e&eb) a mixed system-apparatus state and show that (eb) this effect is fundamentally different from (e) quantum state cloning. Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. Lett. 111, 040401 (2013) DOI:10.1103/PhysRevLett.111.040401 Report number: LA-UR-12-25351 Cite as: arXiv: 1304.2914 [quant-ph] (or arXiv: 1304.2914v2 [quant-ph] for this version) Quantum discord cannot be shared Alexander Streltsov, Wojciech H. Zurek STOP Wojciech H. Zurek, Don N. Page investigate thermodynamic equilibrium of (e) a self-gravitating perfect fluid in (eb) a spherically symmetric system containing (e) a black hole of mass M by means of (e) the Tolman-Oppenheimer-Volkoff (TOV) equation. At r >> 2M its solutions describe (eb) a black-body radiation atmosphere with (e&eb) the Hawking temperature T_BH~1/(8 \pi M) that is increasingly blueshifted as r approaches (eb) 2M. However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments:**

**Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by**

**(e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a firewall-like "Planck cocoon" with entropy close to (e) black hole entropy**

## Prescientiаl Dynamics And Contential Staticity Of Accentuation Attrition Models

## Module Four

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equailibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-445) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Some of the concepts to which model is applicable is as follows:**

**Historically, the first true GUT which was based on (e) the simple Lie group SU(5), was proposed by (e) Howard Georgi and Sheldon Glashow in 1974.[2] The Georgi–Glashow model was preceded by (e) the semisimple Lie algebra Pati–Salam model by (e) Abdus Salam and Jogesh Pati,[3] who pioneered the idea to unify gauge interactions. The acronym GUT was first coined in 1978 by CERN researchers John Ellis, Andrzej Buras, Mary K. Gaillard, and Dimitri Nanopoulos, however in (eb) the final version of their paper [4] they opted for (e) the less anatomical GUM (Grand Unification Mass). Nanopoulos later that year was the first to use [5] the acronym in a paper.[6] The fact that the electric charges of (e) electrons and protons seem to cancel (e&eb) each other exactly to (e) extreme precision is essential for (e) the existence of the macroscopic world as we know it, but this important property of (e) elementary particles is not explained in the Standard Model of particle physics. While the description of strong and weak interactions within (eb) the Standard Model is based on (e) gauge symmetries governed by (e&eb) the simple symmetry groups SU(3) and (e&eb) SU(2) which allow (eb) only discrete charges, the remaining component, the weak hypercharge interaction is described by (e) an abelian symmetry U(1) which in principle allows (eb) for arbitrary charge assignments.[note 1] The observed charge quantization, namely the fact that all known elementary particles carry electric charges which appear to be (=) exact multiples of ⅓ of the "elementary" charge, has led to (eb) the idea that hypercharge interactions and possibly the strong and weak interactions might be embedded in (eb) one Grand Unified interaction described by (e) a single, larger simple symmetry group containing (e) the Standard Model. This would automatically predict (eb) the quantized nature and values of (e) all elementary particle charges. Since this also results in a prediction for (e) the relative strengths of the fundamental interactions which we observe, in particular the weak mixing angle, Grand Unification ideally reduces (e) the number of independent input parameters, but is also constrained by (e) observations. Grand Unification is reminiscent of (e) the unification of electric and (e&eb) magnetic forces by (e) Maxwell's theory of electromagnetism in the 19th century, but its physical implications and mathematical structure are qualitatively different. Physics Letters B Volume 315, Issues 3–4, 7 October 1993, Pages 325-330 Supersymmetric unification without proton decay Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. G. Lazarides, C. Panagiotakopoulos Q. Shafi. Some sentences have been deleted attributable and ascribable to restraints on space. Please excuse me STOP On September 14, 2015 at 09:50:45 UTC the two detectors of the Laser Interferometer Gravitational-Wave Observatory simultaneously observed (e) a transient gravitational-wave signal. The signal sweeps upwards in frequency from 35 to 250 Hz with (e&eb) a peak gravitational-wave strain of (e) 1.0×10−21. It matches the waveform predicted by general relativity for (e) the inspiral and merger of a pair of black holes and the ringdown of the resulting (eb) single black hole. The signal was observed with a**

**matched-filter signal-to-noise ratio of (e) 24 and a false alarm rate estimated to be (=) less than 1 event per 203 000 years, equivalent to (=) a significance greater than 5.1σ. The source lies at (eb) a luminosity distance of 410+160−180 Mpc corresponding to (e&eb) a redshift z=0.09+0.03−0.04. In the source frame, the initial black hole masses are (=) 36+5−4M⊙ and 29+4−4M⊙, and the final black hole mass is (=) 62+4−4M⊙, with (e&eb) 3.0+0.5−0.5M⊙c2 radiated in (eb) gravitational waves. All uncertainties define (eb) 90% credible intervals. These observations demonstrate (eb) the existence of binary stellar-mass black hole systems. Observation of Gravitational Waves from a Binary Black Hole Merger B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 061102 – Published 11 February 2016 This is the first direct detection of gravitational waves and (e&eb) the first observation of a binary black hole merger. Received 21 January 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.061102 STOP This is the second in a series of three papers on (e&eb) systematic analysis of rank 1 Coulomb branch geometries of (e) four dimensional N=2 SCFTs. In the first paper authors developed a strategy for (e) classifying physical rank-1 CB geometries of N=2 SCFTs. They show how to carry out this strategy computationally to construct (eb) the Seiberg-Witten curves and (e&eb) one-forms for all the rank-1 SCFTs. Explicit expressions are given for all cases, with the exception of the Nf=4 SU (2) gauge theory and (e&eb) the En SCFTs which were previously constructed. Classification includes all known rank-1 theories plus (+) a new one with (e&eb) an abelian flavor group, plus (+) nine additional theories whose existence is more speculative. Four of those, reported in first paper, depend on (=) the assumption of new frozen rank-1 SCFTs. Here authors also show that the assumption of the existence of certain rank-0 N=2 SCFTs leads to (eb) five additional consistent rank-1 CB geometries. Subjects: High Energy Physics - Theory (hep-th) Cite as:arXiv: 1601.00011 [hep-th] (or arXiv: 1601.00011v2 [hep-th] for this version Geometric constraints on the space of N=2 SCFTs II: Construction of special Kähler geometries and RG flows Philip C. Argyres, Matteo Lotito, Yongchao Lü, Mario Martone (Submitted on 31 Dec 2015 (v1), last revised 6 Jan 2016 (this version, v2)) STOP Authors propose a model of Dark Energy in which the field currently dominating (e) the energy density of the universe is (=) an "axion field" linearly coupled to (e&eb) the Pontryagin density, tr (F∧F), (i.e., the exterior derivative of (e) the Chern-Simons form) of (e) a massive gauge field. Authors assume that the axion has (e) self-interactions corresponding to (e&eb) a non-trivial (exponential) potential. Authors argue that a non-vanishing magnetic helicity of (e) the gauge field triggers (eb) slow-rolling of the axion at field values far below (eb) the Planck scale. Proposal leads to (eb) a "Tracking Dark Energy Scenario" in which the contribution of the axion energy density to (e) the total energy density is (=) constant (and small) during (e&eb) the early radiation phase, until (e) a secular growth term proportional to (e) the Pontryagin density of (e) the gauge field becomes (=) dominant. The initially small contribution of (e) the axion field to (e) the total energy density is related to the observed small baryon-to-entropy ratio. Some speculations concerning the nature of the gauge field are offered.**

**Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Cite as:arXiv: 1601.00057 [hep-th] (or arXiv:1601.00057v2 [hep-th] for this version) Tracking Dark Energy from Axion-Gauge Field Couplings Stephon Alexander, Robert Brandenberger, Juerg Froehlich (Submitted on 1 Jan 2016 (v1), last revised 23 Jan 2016 (this version, v2)) STOP Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version) Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP In this work authors obtain bounds on (e&eb) the topological Abelian string-vortex and on (e&eb) the string-cigar, by using (e) a new**

**measure of configurational complexity, known as (=) configurational entropy. In this way, the information-theoretical measure of (e) six-dimensional braneworld scenarios is capable to (e) probe situations where (e) the parameters responsible for (e) the brane thickness are arbitrary. The so-called configurational entropy (CE) selects (e&eb) the best value of the parameter in (eb) the model. This is accomplished by (e) minimizing the CE, namely, by (e) selecting the most appropriate parameters in the model that correspond to (e&eb) the most organized system, based upon (e) the Shannon information theory. This information-theoretical measure of complexity provides (eb) a complementary perspective to (e) situations where (e) strictly energy-based arguments are (=) inconclusive. Authors show that the higher the energy the higher the CE, what shows (eb) an important correlation between the energy of the localized field configuration and (e&eb) its associated entropic measure. Subjects: High Energy Physics - Theory (hep-th) DOI:10.1016/j.physletb.2016.02.038 Cite as: arXiv: 1601.00076 [hep-th] (or arXiv: 1601.00076v2 [hep-th] for this version) Bounds on topological Abelian string-vortex and string-cigar from information-entropic measure R. A. C. Correa, D. M. Dantas, C. A. S. Almeida, Roldao da Rocha (Submitted on 1 Jan 2016 (v1), last revised 17 Feb 2016 (this version, v2)) STOP Sum of the Holevo quantity (that bounds the capacity of quantum channels to transmit (e&eb) classical information about an observable) and (e&eb) the quantum discord (a measure of the quantumness of (e) correlations of that observable) yields (eb) an observable-independent total given by (e) the quantum mutual information This split naturally delineates ,dilates and disseminates information about (e&eb) quantum systems accessible to (e) observers -- information that is redundantly transmitted by (e&eb) the environment -- while showing that it is (=) maximized for (e) the quasi-classical pointer observable. Other observables are accessible only via (e&eb) correlations with the pointer observable .Further authors prove an anti-symmetry property relating accessible information and (e&eb) discord. It shows that information becomes (=) objective -- accessible to (e) many observers -- only as quantum information is relegated to (e) correlations with (e&eb) the global environment, (eb) and, therefore, locally inaccessible. The resulting complementarity explains (eb) why, in a quantum Universe, we perceive objective classical reality while (e) flagrantly quantum superpositions are out of reach. Subjects: Quantum Physics (quant-ph) Journal reference: Scientific Reports 3, 1729 (2013) Report number: LAUR-13-26386 Cite as:arXiv: 1303.4659 [quant-ph] (or arXiv: 1303.4659v2 [quant-ph] for this version) Complementarity of quantum discord and classically accessible information Michael Zwolak, Wojciech H. Zurek STOP Niels Bohr proposed that the outcome of the measurement becomes (=) objective and real, and, hence, classical, when (e) its results can be communicated by (e) classical means. In this work authors revisit Bohr's postulate using (e) modern tools from (e) the quantum information theory. They find a full confirmation of Bohr's idea: if a measurement device is in a nonclassical state, (eb) the measurement results cannot be communicated perfectly by classical means. In this case some part of information in the measurement apparatus is lost (e) in the process of communication: the amount of this lost information turns out to be (=) the quantum discord. The information loss occurs even when (e) the apparatus is not entangled with the system of interest. The tools presented in this work allow (eb) to generalize Bohr's postulate: authors show that for pure system-apparatus states (eb) quantum communication does not provide any advantage when (e) measurement results are communicated to (e) more than one recipient. They further demonstrate the superiority of quantum communication to (e&eb) two recipients on (e&eb) a mixed system-apparatus state and show that (eb) this effect is fundamentally different from (e) quantum state cloning. Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. Lett. 111, 040401 (2013) DOI:10.1103/PhysRevLett.111.040401 Report number: LA-UR-12-25351 Cite as: arXiv: 1304.2914 [quant-ph] (or arXiv: 1304.2914v2 [quant-ph] for this version) Quantum discord cannot be shared Alexander Streltsov, Wojciech H. Zurek STOP Wojciech H. Zurek, Don N. Page investigate thermodynamic equilibrium of (e) a self-gravitating perfect fluid in (eb) a spherically symmetric system containing (e) a black hole of mass M by means of (e) the Tolman-Oppenheimer-Volkoff (TOV) equation. At r >> 2M its solutions describe (eb) a black-body radiation atmosphere with (e&eb) the Hawking temperature T_BH~1/(8 \pi M) that is increasingly blueshifted as r approaches (eb) 2M. However, there is no horizon at (eb) the Schwarzschild radius. Instead, the fluid becomes (=) increasingly hot and dense there, piling up into (e&eb) a "firewall" with the peak temperatures and densities reaching (eb) Planck values somewhat below r = 2M. This firewall surrounds (e&eb) a negative point mass residing at r=0, the only singularity of the solution. The entropy of the firewall is comparable to (e&eb) the Bekenstein-Hawking entropy. Comments: Authors study self-gravitating spherically symmetric fluid with (e&eb) a mass of a black hole surrounded by (e) Hawking radiation. Solutions cross r=2M without (e) encountering coordinate singularity to reach a**

**firewall-like "Planck cocoon" with entropy close to (e) black hole entropy. Authors reproduce paper with an updated title and abstract. For a later study with similar results see G. 't Hooft, gr-qc/9706058 Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Physical Review D, Volume 29, Number 4, pp. 628 - 631 (1984) DOI:10.1103/PhysRevD.29.628 Cite as:arXiv:1511.07051 [gr-qc] (or arXiv:1511.07051v1 [gr-qc] for this version) Firewalls, black-hole thermodynamics, and singular solutions of the Tolman-Oppenheimer-Volkoff equation Wojciech H. Zurek, Don N. Page (Submitted on 22 Nov 2015) STOP Frederic Magniez, Ashwin Nayak consider the problem of testing the commutativity of (e) a black-box group specified by its k generators. The complexity (in terms of k) of (e) this problem was first considered by Pak, who gave (eb) a randomized algorithm involving O(k) group operations. They construct a quite optimal quantum algorithm for this problem whose complexity is in O~(k2/3)O~(k2/3). The algorithm uses and highlights the power of the quantization method of Szegedy. For the lower bound of Ω(k2/3)Ω(k2/3), authors give a reduction from (e) a special case of Element Distinctness to our problem. Along the way, they prove the optimality of the algorithm of (e) Pak for the randomized model. Algorithmica July 2007, Volume 48, Issue 3, pp 221-232 First online: 25 June 2007 Quantum Complexity of Testing Group Commutativity Frederic Magniez, Ashwin Nayak STOP Pawel Wocjan, Jon Yard analyze relationships between quantum computation and (e&eb) a family of generalizations of the Jones polynomial. Extending recent work by Aharonov et al., authors give efficient quantum circuits for (e) implementing the unitary Jones-Wenzl representations of the braid group. We use these to provide new quantum algorithms for (e) approximately evaluating a family of specializations of the HOMFLYPT two-variable polynomial of trace closures of braids. They also give (eb) algorithms for approximating the Jones polynomial of a general class of closures of (e) braids at roots of unity. Next authors provide (eb) a self-contained proof of a result of (e) Freedman et al. that any quantum computation can be replaced by (e&eb) an additive approximation of the Jones polynomial, evaluated at (eb) almost any primitive root of unity. Proof encodes two-qubit unitaries into (e&eb) the rectangular representation of (e) the eight-strand braid group. Authors then give QCMA-complete and (e&eb) PSPACE-complete problems which are based on braids. They conclude with direct proofs that evaluating the Jones polynomial of the plat closure at most primitive roots of unity is a #P-hard problem, while learning its most significant bit is PP-hard, circumventing (e) the usual route through the Tutte polynomial and graph coloring. Subjects: Quantum Physics (quant-ph); Geometric Topology (math.GT) Cite as: arXiv:quant-ph/0603069 (or arXiv:quant-ph/0603069v3 for this version) The Jones polynomial: quantum algorithms and applications in quantum complexity theory Pawel Wocjan, Jon Yard STOP Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan address the task of switching off the Hamiltonian of (e) a system by removing (e) all internal and system-environment couplings. Authors propose dynamical decoupling schemes, which use (e) only bounded-strength controls, for quantum many-body systems with local system Hamiltonians and local environmental couplings. To do so, they introduce (eb) the combinatorial concept of (e) balanced-cycle orthogonal arrays (BOAs) and show how to construct them from (e) classical error-correcting codes. The derived decoupling schemes may be useful as (=) a primitive for more complex schemes, e.g., for Hamiltonian simulation. For the case of n qubits and a 2-local Hamiltonian, the length of the resulting decoupling (e&eb) scheme scales as O (nlogn), improving over (e&eb) the previously best-known schemes that scaled quadratically with (e&eb) n. More generally, using balanced-cycle orthogonal arrays constructed from families of BCH codes, we show that bounded-strength decoupling for any ℓ-local Hamiltonian, where ℓ≥2, can be achieved using decoupling schemes of length at most O(nℓ−1logn). Subjects:Quantum Physics (quant-ph); Emerging Technologies (cs.ET) Journal reference: IEEE Trans. Inf. Theory, 62(5): 2881-2894, 2016 DOI: 10.1109/TIT.2016.2535183 Report number: MIT-CTP /4705 Cite as: arXiv: 1509.00408 [quant-ph] (or arXiv: 1509.00408v2 [quant-ph] for this version) Improved bounded-strength decoupling schemes for local Hamiltonians Adam D. Bookatz, Martin Roetteler, and Pawel Wocjan STOP Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola propose dynamical control schemes for Hamiltonian simulation in many-body quantum systems that avoid instantaneous control operations and rely solely on (e) realistic bounded-strength control Hamiltonians. Each simulation protocol consists of (e) periodic repetitions of a basic control block, constructed as (=) a suitable modification of an "Eulerian decoupling cycle," that would otherwise implement (eb) a trivial (zero) target Hamiltonian. For an open quantum system coupled to (e) an uncontrollable environment, approach may be employed to (e) engineer an effective evolution that simulates (e&eb) a target Hamiltonian on the system, while suppressing (e) unwanted decoherence to the leading order. Authors present illustrative applications to (e) both closed- and open-system simulation settings, with (e&eb)**

**emphasis on simulation of non-local (two-body) Hamiltonians using (e) only local (one-body) controls. In particular, authors provide (eb) simulation schemes applicable to (e&eb) Heisenberg-coupled spin chains exposed to (e) general linear decoherence, and show how to simulate (e&eb) Kitaev's honeycomb lattice Hamiltonian starting from Ising-coupled qubits, as (=) potentially relevant to (e&eb) the dynamical generation of a topologically protected (e&eb) quantum memory. Additional implications for quantum information processing are discussed. Subjects: Quantum Physics (quant-ph) Journal reference: New J. Phys. 16, 045021 (2014) DOI: 10.1088/1367-2630/16/4/045021 Cite as:arXiv: 1310.4153 [quant-ph]**

**(or arXiv: 1310.4153v1 [quant-ph] for this version) Hamiltonian quantum simulation with bounded-strength controls Adam D. Bookatz, Pawel Wocjan, and Lorenza Viola STOP Felix Fontein, Pawel Wocjan present a quantum algorithm for computing the period lattice of (e) infrastructures of fixed dimension. The algorithm applies to (e&eb) infrastructures that satisfy certain conditions. The latter are always fulfilled for (e) infrastructures obtained from (e) global fields, i.e., algebraic number fields and function fields with (e&eb) finite constant fields. The first main contributions is an exponentially better method for (e) sampling approximations of vectors of the dual lattice of (e) the period lattice than the methods outlined in the works of (e) Hallgren and Schmidt and Vollmer. This new method improves (eb) the success probability by (e) a factor of at least 2^ {n^2-1} where (e) n is the dimension. The second main contribution is a rigorous and complete proof that the running time of the algorithm is (=) polynomial in the logarithm of (e) the determinant of the period lattice and (e&eb) exponential in n. The third contribution is the determination of (e) an explicit lower bound on (e&eb) the success probability of our algorithm which greatly improves on (e&eb) the bounds given in the above works. The exponential scaling seems inevitable because (e) the best currently known methods for (e) carrying out fundamental arithmetic operations in infrastructures obtained from (e) algebraic number fields take exponential time. In contrast, the problem of computing (e&eb) the period lattice of infrastructures arising from function fields can be solved without (e) the exponential dependence on (e&eb) the dimension n since this problem reduces efficiently to the abelian hidden subgroup problem. This is also true for (e) other important computational problems in (eb) algebraic geometry. The running time of the best classical algorithms for (e) infrastructures arising from (e) global fields increases subexponentially with (e&eb) the determinant of the period lattice. Subjects: Quantum Physics (quant-ph) Cite as: arXiv: 1111.1348 [quant-ph]**

**(or arXiv: 1111.1348v2 [quant-ph] for this version) Quantum Algorithm for Computing the Period Lattice of an Infrastructure Felix Fontein, Pawel Wocjan STOP Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan advocate a new approach of (e) addressing hidden structure problems and (e&eb) finding efficient quantum algorithms. Authors introduce and investigate the Hidden Symmetry Subgroup Problem (HSSP), which is a generalization of the well-studied Hidden Subgroup Problem (HSP). Given a group acting on (e&eb) a set and an oracle whose level sets define (eb) a partition of the set; the task is to recover (e) the subgroup of symmetries of (e) this partition inside the group. The HSSP provides (eb) a unifying framework that, besides the HSP, encompasses (e) a wide range of algebraic oracle problems, including (e) quadratic hidden polynomial problems. While the HSSP can have (e) provably exponential quantum query complexity, authors obtain efficient quantum algorithms for various interesting cases. To achieve this, authors present a general method for (e) reducing the HSSP to the HSP, which works efficiently in several cases related to (e&eb) symmetries of polynomials. The HSSP therefore connects in a rather surprising way certain hidden polynomial problems with (e&eb) the HSP. Using this connection, authors obtain the first efficient quantum algorithm for (e) the hidden polynomial problem for (e) multivariate quadratic polynomials over fields of (e) constant characteristic. Authors also apply the new methods to (e&eb) polynomial function graph problems and (e&eb) present an efficient quantum procedure for (e) constant degree multivariate polynomials over (e&eb) any field. This result improves in several ways the currently known algorithms. Subjects: Quantum Physics (quant-ph) Journal reference: SIAM Journal on Computing 42:5 (2013) DOI: 10.1137/120864416 Cite as: arXiv: 1107.2189 [quant-ph]**

**(or arXiv: 1107.2189v2 [quant-ph] for this version) Hidden Symmetry Subgroup Problems Thomas Decker, Gábor Ivanyos, Miklos Santha, Pawel Wocjan STOP Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan present an efficient general method for realizing a quantum walk operator corresponding to (e&eb) an arbitrary sparse classical random walk. Approach is based on Grover and Rudolph's method for preparing coherent versions of (e) efficiently integrable probability distributions. This method is intended for (e) use in quantum walk algorithms with polynomial speedups, whose complexity is usually measured in terms of (e&eb) how many times we have to apply a step of a quantum walk, compared to (e&eb) the number of necessary classical Markov chain steps. Authors consider a finer notion of complexity including (e) the**

**number of elementary gates it takes to implement each step of the quantum walk with (e&eb) some desired accuracy. The difference in complexity for (e) various implementation approaches is that method scales linearly in the sparsity parameter and (e&eb) poly-logarithmically with (e&eb) the inverse of the desired precision. Best previously known general methods either scale quadratically in (eb) the sparsity parameter, or polynomially in (eb) the inverse precision. Approach is especially relevant for (e) implementing quantum walks corresponding to (e&eb) classical random walks like those used in (eb) the classical algorithms for approximating permanents and sampling from binary contingency tables. In those algorithms, the sparsity parameter grows with (e&eb) the problem size, while (e) maintaining high precision is required. Subjects: Quantum Physics (quant-ph) Journal reference: Quantum Information and Computation, Vol.10, No.5&6, pp0420-0434 (2010) Cite as: arXiv: 0903.3465 [quant-ph] (or arXiv: 0903.3465v3 [quant-ph] for this version) Efficient Circuits for Quantum Walks Chen-Fu Chiang, Daniel Nagaj, Pawel Wocjan STOP Preparing the ground state of a system of interacting classical particles is (=) an NP-hard problem. Thus, there is in general no better algorithm to (e) solve this problem than exhaustively going through all N configurations of (e) the system to determine (eb) the one with lowest energy, requiring (e) a running time proportional to N. A quantum computer, if it could be built, could solve (e) this problem in time sqrt (N). Here, authors present a powerful extension of this result to (e) the case of interacting quantum particles, demonstrating (eb)that a quantum computer can prepare the ground state of (e) a quantum system as efficiently as it does for (e) classical systems. Subjects: Quantum Physics (quant-ph) Journal reference: Phys.Rev.Lett.102:130503, 2009 DOI:10.1103/PhysRevLett.102.130503 Cite as: arXiv: 0809.2705 [quant-ph] (or arXiv: 0809.2705v1 [quant-ph] for this version) Preparing ground states of quantum many-body systems on a quantum computer David Poulin, Pawel Wocjan STOP Daniel Nagaj, Pawel Wocjan construct a simple translationally invariant, nearest-neighbor Hamiltonian on (e&eb) a chain of 10-dimensional qudits that makes it possible to (e) realize universal quantum computing without (e) any external control during (e&eb) the computational process. Authors only require the ability to prepare an initial computational basis state which encodes (e) both the quantum circuit and its input. The computational process is (=) then carried out by (e) the autonomous Hamiltonian time evolution. After a time polynomially long in the size of (e) the quantum circuit has passed, the result of (e) the computation is obtained with (e&eb) high probability by (e) measuring a few qudits in the computational basis. This result also implies (eb) that there cannot (e) exist efficient classical simulation methods for (e) generic translationally invariant nearest-neighbor Hamiltonians on (e&eb) qudit chains, unless (e) quantum computers can be efficiently simulated by (e&eb) classical computers (or, put in complexity theoretic terms, unless BPP=BQP). Subjects: Quantum Physics (quant-ph) Journal reference: Phys. Rev. A 78, 032311 (2008) DOI: 10.1103/PhysRevA.78.032311 Cite as: arXiv: 0802.0886 [quant-ph] (or arXiv: 0802.0886v4 [quant-ph] for this version) Hamiltonian Quantum Cellular Automata in 1D Daniel Nagaj, Pawel Wocjan**

## 64. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_4}{dt} = -b_4 Q_4 + g_4(g_4)Q_4$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_4(0) = 0, \quad g_4(P_4) > 0, \ P_4 \ 0, \qquad \lim_{P_4 \to \infty} g_4(P_4) = (K_4)_4 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_4 + (\rho_a)_4 + D_4(a_4, P_4, Q_4)\rho_4 + 0, \qquad 0 < a_4 \ \infty,$$

$$\rho_4(t, 0) = \int_0^\infty n_4(a_4, P_4, Q_4)\rho_4(t_4, a_4)da_4$$

Where $D_4$ is the mortality function and $n_4$ is the fertility function.

Put $_4$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_4}{dt}\right]_4 = x_4\left[\frac{dQ_2}{dt}\right]_4 - x_2\left[\frac{dQ_4}{dt}\right]_4$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_4 - x_4\left[\frac{dQ_2}{dt}\right]_4$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_4 - x_6\left[\frac{dQ_3}{dt}\right]_4$$

Where $x_4, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_4$ and $g_4(P_4)Q_4$

Now when you add the term $(b_4Q_4)$ with $g_4(P_4)(Q_4)$ we are infect see is that $(b_4Q_4)$ is credited incrementally part of $g_4(P_4)(Q_4)$ or instant nearly, when subtracted from $\frac{dQ_4}{dt}$, it means that $g_4(P_4)(Q_4)$ is slowly eaten by $\frac{dQ_4}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_4 = \mu_4(a_4) + (f_4)_4(P_4)_4 + (f_2)_4(Q)_4 \quad (2)$$

$$n_4 = \beta_4(a_4)h_4(Q_4), \quad (3)$$

Where $\mu_4$the age- dependent debit rate whit is $\mu_4(a_4)$ $\forall a > 0$; $(f_4)_4$ is the density- dependent death rate with $f_4(0) = 0$, $(f_4^4)_4(P_4)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_4)_4(P_4) = (k_4)_4$ ; $f_2$ is the functional response of creditors to prey with $[f_2]_4(0) = 0, (f_2^4)_4(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_4(Q_4) = [k_2]$ ; $\beta$ is the age dependent birth rate with $\beta(a)$ $0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 4, h^4(Q)$ $0, \forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_4 + (\rho_a)_4 = -\mu_4(a_4) + (f_4)_4P_4) + \big(f_2)_4(Q_4)\big)(\rho_4)_4 \quad (4)$$

$$\rho_4(t_4, 0) = h_4(Q_4)\int_0^\infty \beta_4(a_4)\rho_4(t, a_4)da, \quad (5)$$

$$\frac{dQ_4}{dt}$$

$$= -b_4Q_4 + g_4(P_4)Q_4 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_4$ and $Q_4$ of (4-6) satisfy

$$0 \leq \rho_4(t, a_4) \leq \rho_4^*(t, a_4),$$

$$0 \leq Q_4(t) \leq Q_4^*(t),$$

Where $\rho_4^*(t, a_4)$ and $Q_4^*(t)$, are solutions of

$$(\rho_t^*)_4 + (\rho_a^*)_4 = -(\mu_4(a_4) + (f_4)_4(P_4^*) + (f_2)_4 Q_4^*))\, \rho_4^*, \quad (7)$$

$$\rho_4^*(t, 0) = \int_0^\infty \beta_4(a_4)\, \rho_4^*(t, a_4) da_4 \quad (8)$$

$$\frac{dQ_4^*}{dt} = -b_4 Q_4^* + g_4(P_4^*)Q_4^*, \quad (9)$$

With $P_4^*(t) = \int_0^\infty \rho_4^*(t, a_4) da_4$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_4^*(t)}{dt} = P_4^*(t)\big(r_4 - (f_4)(P_4^*) - (f_2)_4(Q^*)\big), \quad (10)$$

$$\frac{dQ_4^*(t)}{dt} = Q_4^*(t)(-b_4 + g_4(P_4^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_4(a) e^{-\int_0^a \mu_4(\tau)dt} e^{-ra}\, da_4 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_4^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^{\infty} \beta_4(a_4) e^{-\int_0^a \mu_4(t)dt}\, da_4 > 1, \qquad (H1)$$

Which guarantees $r_4 > 0$.

*Remark.* If $\beta_4(a) = [\beta_0]_4 a_4 e^{-\alpha_4 a_4}, \mu_4(a_4) \equiv \mu_4$ then the solution of (12) is

$$r_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_4]_4(P_4) = 0, [f_2](Q_4) = \lambda_4 Q_4\ h_4(Q_4) = 1, \qquad and\ \ g_4(P_4) = c_4 P_4$$

$$[\rho_4]_4 + [\rho_a]_4 = -(\mu_4 + \lambda_4 Q_4)\rho_4$$

$$\rho_4(t,0) = \int_0^{\infty} [\beta_0]_4 a_4 e^{-\alpha_4 a_4} \rho_4(t, a_4)\, da_4$$

$$\frac{dQ_4}{dt} = -b_4 Q_4 + c_4 P_4 Q_4$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_4}{dt} = \omega_4 P_4 - \lambda_4 P_4 Q_4$$

$$\frac{dQ_4}{dt} = b_4 Q_4 - c_4 P_4 Q_4$$

Where $\omega_4 = \sqrt{\beta_4} - \alpha_4 - \mu_4$

## 65. THE EQUILIBRIUM $(\rho_0)_4(a_4), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_4(a_4), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_4 + (\rho_a)_4 = -\big(\mu_4(a_4) + (f_4)_4(P_4)\big)\rho,$$

$$\rho_4(t,0) = \int_0^\infty \beta_4(a_4)\rho_4(t,a_4)da_4 \qquad (13)$$

The solution of (13) has the form of

$$\rho_4(t,a_4) = Ce_4^{r(t-a_4)} - \int_0^a (\tau)d\tau - \int_0^4 (f_4)_4\big(P_4(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_4(t) = C_4 e^{r_4 t - \int_0^t (f_4)_4(P_4(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_4^{-4}(r)$ if $k_4 > r$. This is equivalent to

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{(k_4)_4 a_4} da_4 < 1.$$

In fact, if we define

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau)d\tau} e^{-s_4 a_4}\, da_4$$

$R(s)$ is a decreasing function. $R_4(r_4) = 1$ And $R_4(k_4)_4 < 1$ imply $r_4 < (k_4)_4$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_4(a_4), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{k_4 a_4} da_4$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_4 \equiv \int_0^\infty (\rho_0)_{44}(a_4)da_4$ *satisfies*

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau)d\tau} e^{-(f_4)_4(P_0)a_4} da$$

$$= 1 \qquad (11)$$

Next, we study the stability of $(\rho_0)_4(a_4), 0)$

Let $u_4(t, a_4) = \rho(t, a_4) - \rho_0(a_4)$ Then

$$N_4(t) \equiv \int_0^\infty u_4(t, a_4)da_4 = \int_0^\infty \rho_4(t, a_4)da_4 - \int_0^\infty (\rho_0)_4(a_4)da = P_4(t) - P_0)_4$$

And

$$(u_4)_4 + (u_a)_4 = (\rho_t)_4 + (\rho_a)_4 - (\rho_0^4)_4(a_4)$$

$$= -\mu_4(a_4) + (f_4)_4(P_4) + (f_2)_4(Q_4))\,(u_4 + (\rho_0)_4) + \mu_4(a_4)(\rho_0)_4 + (f_4)_4(P_0)_4(\rho_0)_4$$

$$= -\mu_4(a_4)u_4 - (f_4)_4(P_4) + (f_2)_4(Q_4))\,u_4 - (\rho_0)_4\,(f_4)_4(P_4) - (f_4)_4(P_0)_4$$

$$+ (f_2)_4(Q_4)), \qquad (15)$$

Where $(\rho_0^4)_4(a_4)$ denotes $d(\rho_0)_4(a_4)/da_4$

Expansion $(f_4)_4(P_4)$ and $(f_2)_4(Q_4)$ at $(P_0)_4$ and 0 yield respectively

$$(f_4)_4(P_4) = (f_4)_4((P_0)_4 + N_4)$$

$$= (f_4)_4(P_0)_4 + (f_4^4)_4(P_0)_4 N_4 + o_4(N_4), \quad (16)$$

$$(f_2)_4(Q_4)$$

$$= (f_2^4)_4(0)Q + o_4(Q_4) \quad (17)$$

Write $f_4(P_0) = f_{40}, f_4'(P_0) = f_{40}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_4)_4 + (u_a)_4 \approx -(\mu_4(a_4) + (f_{40})_4)u - (f_{40}^4)_4(\rho_0)_4(N_4) - (f_{20}^4)_4(\rho_0)_4(Q_4)$$

$$u_4(t,0) = \rho_4(t,0)(\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a) + (\rho_0)_4(a_4)da_4 - (\rho_0)_4(0)$$

$$= h_4(Q_4)\int_0^\infty \beta_4(a_4)(u_4(t,a_4)da_4 + (h_4(Q_4) - 4)(\rho_0)_4(0) \quad (18)$$

Expanding

$$h_4(Q_4) = 1 + (h^4)_4(0)Q_4 + 0(Q_4) \equiv 1 + h_0'Q_4 + 0Q_4$$

And substituting it into (48), we arrive at

$$u_4(t,0) \approx \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + (h_{a_4}')(\rho_0)_4 Q_4$$

The linearization is the following :

$$[u_4] + [u_4] = -(\mu_4(a_4) + [t_{40}]_4 u_4 - [\alpha_4]_4(a_4)N_4 - [\alpha_2]_4(a_4)Q_4$$

$$u_4(t,0) = \int_0^\infty \beta_4(a_4)u_4(t,a_4)da_4 + [\xi_4]Q_4 \quad (19)$$

$$\frac{dQ_4}{dt} = -b_4 Q_4$$

where

$$[\alpha_4]_4(a_4) = [t_4^4 a]_4[p, o]_4(a_4) d2(a_4) = [f_{20}^4]_4(\rho\, o\,)_4(a_4)\xi = [\mathrm{h\ o}^4]_4[\rho \mathrm{o}]_4(\mathrm{o}) \text{ and } \mathrm{b}_4 = \mathrm{b}_4 - g_4(P_4)$$

If $g_4(P_0) > b$, clearly, $(\rho_0)_4(a_4), 0)$ is unstable and the predator population is persistent. Assume $g_4(P_0)_4 < b$.

Then, $Q_4(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_4 + u_a = -(\mu_4(a_4) + (f_{40})_4)u_4 - (\alpha_4)_4(a_4)N_4$$

$$= -(\mu_4(a_4) + f_{40})u_4 - \alpha_4(a_4)\int_0^\infty u_4(t, a_4)da_4$$

$$u_4(t, 0) = \int_0^\infty \beta_4\,(a_4)u_4(t, a_4)da_4$$

Using the method developed by Cushing [2, 3], we suppose

$$u_4(t, a_4) = \tilde{\mathrm{u}}_4(a_4)^{e^{\gamma(t-a_4)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_4}\, da_4$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_4)$ gives,

$$\tilde{\mathrm{u}}_4(a_4) = e^{-\int_0^a \mu_4(\tau_4)d\tau - (f_{40})_4 a}\left(\tilde{\mathrm{u}}_4(0) - w\int_0^a e^{\int_0^\tau \mu_4(s_4)ds + (f_{40})}\tau_4\,(\alpha_4)_4(\tau_4)e^{\gamma_4\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1-\int_0^{\infty}\beta_4(a_4)e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\right)\tilde{u}_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a(\alpha_4)_4(\tau_4)e^{-\gamma_4\tau_4}e^{(f_{40})_4\tau_4+\int_0^\tau \mu_4(s_4)ds}d\tau_4\,da_4$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}da\right)(\tilde{u}_4(0))_4$$

$$+\left(-4-\int_0^{\infty}e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(\tau_4)d\tau}\int_0^a e^{(f_{40})_4\tau_4+\int_0^\tau \mu_4(s_4)ds+y\tau}\alpha_4(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_4 a_4-(f_{40})_4 a_4-\int_0^a \mu_4(t_4)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_4\,(a_4)Y_4(\gamma_4,a_4)da_4\right)u_4(0)$$

$$+w\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)\int_0^a(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4=0,$$

(24)

$$\left(\int_0^{\infty}Y_4(\gamma_4,a_4)da_4\right)\tilde{u}_4(0)$$

$$+\left(-1-\int_0^{\infty} Y_4(\gamma_4,a_4)\int_0^{a}(\alpha_4)_4(\tau)_4(Y^{-4})_4(\gamma_4,\tau_4)d\tau_4\,da_4\right)w=0.$$

Equqtion (20) has a solution of the form of $u_4(t,a_4)=u_4(a_4)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_4(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_4(\gamma_4,a_4)=\int_0^a(\rho_0)_4(\tau_4)Y^{-4}(\gamma_4,a_4)d\tau_4$. We obtain the characteristic equation

$$f_{40}^{l}{}_4\left(\int_0^{\infty} Y_4(\gamma_4,a_4)da_4\right)\left(\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$+\left(1-\int_0^{\infty}\beta_4(a_4)Y_4(\gamma_4,a_4)da\right)\left(1+(f_{40}^{l})_4\int_0^{\infty}Y_4(\gamma_4,a_4)I_4(\gamma_4,a_4)da_4\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_4(a_4)=(\rho_0)_4(0)e^{-\int_0^a\mu_4(\tau_4)d\tau_4-(f_{40})_4a_4}$$

$$I_4(\gamma_4,a_4)=\int_0^a(\rho_0)_4(\tau_4)e^{\gamma_4\tau_4+(f_{40})_4\tau+\int_0^{\tau}\mu_4(s_4)ds_4}d\tau$$

$$=\frac{(\rho_0)_4(0)}{\gamma_4}(e^{\gamma_4a_4}-1)$$

Let

$$E_4(a_4)=e^{-(f_{40})_4a_4-\int_0^a\mu_4(\tau_4)d\tau_4}$$

Then $[Y_4]_4[\gamma_4]_4, a_4) = E_4(a_4)e^{-\gamma_4 a_4}$. by a simple calculation,

$$\int_0^\infty \beta_4(a_4)Y_4(\gamma_4, a_4)I(\gamma_4, a_4)da = \frac{[\rho_0]_4(0)}{\gamma}\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}\ da\right);$$

$$\int_0^\infty Y_4(\gamma_4, a_4)I(\gamma_4, a_4)da = \frac{[\rho_0]_4(0)}{\gamma_4}\left(\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}\ da_4\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_4(0)}{\gamma_4}[f_{40}^4]\int_0^\infty E_4(a_4)e^{-\gamma_4 a_4}da_4\ (\int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4)$$

$$+\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right).$$

$$\left(1 + \frac{(\rho_0)_4(0)}{\gamma_4}(f_{40}^{'})_4\left(\int_0^\infty E_4(\,a_4)da_4 - \int_0^\infty E_4(\,a_4)e^{-\gamma_4 a_4}da_4\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4\right)(1 + \frac{(f_{40}^{'})_4}{\gamma_4}(\rho_0)_4(0)\int_0^\infty E_4(\,a_4)da_4) = 0.$$

We need ũ(a) $\not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1 - \int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)da_4 = 1 - \int_0^\infty \beta_4(a_4)E_4(a_4)e^{-\gamma_4 a_4}da_4 \neq 0,$$

Since $\int_0^\infty \beta_4(a_4)Y_4(\gamma_4 a_4)\ \int_0^a \alpha_4\,(\tau_4)Y^{-4}(\gamma_4, \tau_4)d\tau_4 da_4 > 0$. Therefore, the characteristic solution is

$$\gamma_4 = -(f_{40}^{'})_4(\rho_0)_4 \int_0^{\infty} E_4(a_4)da_4.$$

Since $f_4'(P_0)_4 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_4), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_4), 0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_4 + (\mu_4(a_4) + [f_4]_4(P_4^*) + [f_2]_4(Q_4^*))P_4^* = 0, \quad (26)$$

$$p_4^*(0) = h_4(Q_4^*)\int_0^{\infty} \beta_4(a_4)p_4^*(a_J)p_4^*(a_4)da_4 \quad (27)$$

$$g_4(P_4^*) - b_4 = 0, \quad (28)$$

Where $P_4^* = \int_0^{\infty} p_4^*(a_4)\, da_4$

Equation (28) yields $P^* = g^{-4}(b)$,

From (26),

$p_4^*(a) = p_4^*(0)e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}$

Substituting into (27), we need to solve

$$\left(1 - h_4(Q_4^*)\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(Q_4^*)a_4}da_4\right)p_4^*(0) = 0$$

For $Q_4^*$

Fix $p_4^*(0)$ and $P_4^*$ Let

$$H_4(z_4) = 1 - h_4(z_4)\int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}da_{,4}$$

And assume that

$$[k_3]_4\int_0^\infty \beta_4(a_4)e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(g_4^{-4}(b_4))a_4 - [k_2]_4 a_4}\, da_4 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_4(z_4)] > 0.$$

Clearly,

$$\frac{dH_4}{dz_4} = h_4(z)(\,[f_2']_4(z_4)\,)\int_0^\infty a_4\beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - f_4(P_4^*)a_4 - f_2(z_4)a_4}da_{\square}$$

$-h'(z_4)\int_0^\infty \beta_4(a_4)\, e^{\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4 - [f_2]_4(z_4)a_4}\quad da_4$

the assumptions on $\left[f_2^4\right]_4 and\ [h^4]^4 dH_4/dz > 0$. Hence, if

$$H_4(0) = 1 - \int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4(P_4^*)a_4} da_4 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_4(0) < 0$ follows if $g_4^-(b_4) < [f_4^{-4}]_4(r_4)$, where $r_4$ is the solution of (12).

Indeed, since

$$R_4(s_4) = \int_0^\infty \beta_4(a_4)\, e^{-\int_0^a \mu_4(\tau_4)d\tau_4 e^{-y\frac{a}{4}} da_4}$$

Is a decreasing function of $s_4$ when $P_4^* = [g^{-4}]_4(b_4) < [f_4^{-4}]_4(r_4)$, that is $[f_4]_4(P_4^*) < r$, $R\left(f_4(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_4(g^{-4}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4)d\tau_4 - [f_4]_4\left(g_4^{-4}(b)\right)a_4} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_4^*(a_4), Q_4^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_4^*(a_4), Q_4^*)$ is more complicated. We again consider perturbations $m_4(t_4 a_4) = p_4(t_4, a_4) - p_4^*(a_4)$ and $W_4(t) = Q_4(t) - Q^*$.Then

$$M_4(t) \equiv \int_0^\infty m_4(t, a_4) da_4 = P_4(t) - P_4^*,$$

$$[m_4]_4 + [m_a]_4 = [p_4]_4 + [p_a]_4 - [p^*]_4(a_4) \approx -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4$$

$$-[f_{4*}^4]_4 p_4^*(a_4)M - [f_{2*}^4]_4 p_4^*(\mathrm{a})W_4,$$

$$m_4(t,0)p_4(t,0) - p_4^*(0) \approx h_* \int_0^\infty \beta_4(a_4)m_4(t,a_4)\,da + \frac{[h_*^4]_4}{[h_*]_4}\, p_4^*(0)W_4$$

$$\frac{dW_4}{dt} = \frac{dQ_4}{dt} \approx [g_*^4]_4 Q_4^* M_{,4}$$

Where $[f_{4*}]_4$ denoting $[f_4(P^*)]_4$, $[f_{4*}]_4$ denoting $[f_2(P^*)]_4$, $[f_{4*}^{'}]_4$ denoting $[f^{'}{}_4(P^*)]_4[f_{2*}^{'}]_4$ denoting $[f_2]_4$ $(Q_4^*,\quad h_{*}^{}$ denoting $h_*^{'}$ denoting $h_4^{'}$ $(Q_4^*)$, and $[g_*^{'}]_4$

Denoting $[g^{'}]_4(P_4^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_4]_4 + [m_a]_4 = -(\mu_4(a_4) + [f_{4*}]_4 + [f_{2*}]_4)m_4 - [f^{'}{}_{4*}]_4[p_4^*](a_4)M_4 - [f_{2*}^{'}]_4 p_4^*(a_4)W_4$$

$$m_4(t,0) = [h_*]_4 \int_0^\infty \beta_4(a_4)m_4(t,a_4)da_4 + \frac{[h_*^4]_4}{[h^*]_4}[p^*]_4(0)\,W_{,4} \qquad (30)$$

$$\frac{dW_4}{dt} = [g_*^4]_4[Q_4^*]M_4$$

This is equation is important in the following cases: problem of using observed correlations to infer causal relations is relevant to (e) a wide variety of scientific disciplines. Yet given correlations between just two classical variables, (eb) it is impossible to determine whether they arose from (e) a causal influence of one on the other or a common cause influencing both, unless one can implement (eb) a randomized intervention. **Katja Ried, Megan Agnew,**

**Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. Resch** here consider the problem of causal inference for (e) quantum variables. Authors introduce causal tomography, which unifies and generalizes (e&eb) conventional quantum tomography schemes to provide (eb) a complete solution to the causal inference problem using (e) a quantum analogue of a randomized trial. Authors furthermore show that, in contrast to the classical case, observed quantum correlations alone can sometimes provide (eb) a solution. They implement a quantum-optical experiment that allows (eb) to control (e&eb) the causal relation between two optical modes, and two measurement schemes -- one with (e&eb) and one without (e) randomization -- that extract this relation from (e) the observed correlations. Results show that entanglement and (e&eb) coherence, known to be central to (e) quantum information processing, also provide (eb) a quantum advantage for (e) causal inference. Subjects: Quantum Physics (quant-ph); Learning (cs.LG); General Relativity and Quantum Cosmology (gr-qc); Machine Learning (stat.ML) Journal reference: Nat Phys 11, 414-420 (2015) DOI: 10.1038/nphys3266 Cite as: arXiv: 1406.5036 [quant-ph] (or arXiv: 1406.5036v1 [quant-ph] for this version) **Inferring causal structure: a quantum advantage Katja Ried, Megan Agnew, Lydia Vermeyden, Dominik Janzing, Robert W. Spekkens, Kevin J. ReschSTOP D. T. Son, A. O. Starinets** review recent progress in applying the AdS/CFT correspondence to (e&eb) finite-temperature field theory. In particular, they show how the hydrodynamic behavior of field theory is reflected in (e&eb) the low-momentum limit of (e) correlation functions computed through (e&eb) a real-time AdS/CFT prescription, which they formulate. Authors also show how the hydrodynamic modes in field theory correspond to (e&eb) the low-lying quasinormal modes of (e) the AdS black p-brane metric. They provide a proof of the universality of (e) the viscosity/entropy ratio within (eb) a class of theories with (e&eb) gravity duals and formulate (eb) a viscosity bound conjecture. Possible implications for real systems are mentioned. Subjects: High Energy Physics - Theory (hep-th) Journal reference: Ann.Rev.Nucl.Part.Sci.57:95-118,2007 DOI: 10.1146/annurev.nucl.57.090506.123120Report number: INT PUB 07-02Cite as: arXiv:0704.0240 [hep-th]

(or arXiv:0704.0240v2 [hep-th] for this version) An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society **Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP** Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv:

1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version **Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016) STOP Viscosity, Black Holes, and Quantum Field Theory D. T. Son, A. O. Starinets STOP Authors** study information retrieval from evaporating black holes, assuming that the internal dynamics of a black hole is unitary and rapidly mixing, and assuming that the retriever has unlimited control over the emitted Hawking radiation. If the evaporation of the black hole has already proceeded past the ``half-way" point, where (e) half of the initial entropy has been radiated away, then (eb) additional quantum information deposited in (eb) the black hole is revealed in (eb) the Hawking radiation very rapidly. Information deposited prior to the half-way point remains concealed until (e) the half-way point, and then emerges (eb) quickly. These conclusions hold because (e) typical local quantum circuits are efficient encoders for (e) quantum error-correcting codes that nearly achieve (eb) the capacity of the quantum erasure channel. Estimate of a black hole's information retention time, based on (e) speculative dynamical assumptions, is just barely compatible with (e&eb) the black hole complementarity hypothesis. **Black holes as mirrors: quantum information in random subsystems Patrick Hayden1 and John Preskill2 Published 26 September 2007 Journal of High Energy Physics, Volume 2007, JHEP09 (2007) STOP** By analyzing the key properties of black holes from the point of (e) view of quantum information, authors derive a model-independent picture of (e) black hole quantum computing. It has been noticed that this picture exhibits (eb) striking similarities with (e&eb) quantum critical condensates, allowing (eb) the use of a common language to describe (eb) quantum computing in both systems. Authors analyze such quantum computing by allowing coupling to (e&eb) external modes, under the condition that the external influence must be (=) soft-enough in order not to (e) offset the basic properties of the system. Authors derive model-independent bounds on some crucial time-scales, such as (=) the times of gate operation, decoherence, maximal entanglement and total scrambling. They show that for black hole type quantum computers all these time-scales are of (e) the order of the

black hole half-life time. Furthermore, authors construct explicitly a set of Hamiltonians that generates (eb) a universal set of quantum gates for (e) the black hole type computer. They find that the gates work at maximal energy efficiency. Furthermore, authors establish a fundamental bound on the complexity of quantum circuits encoded on (eb&e) these systems, and characterize the unitary operations that are (=) implementable. It becomes apparent that the computational power is very limited due to (e) the fact that the black hole life-time is of the same order of (e) the gate operation time. As a consequence, it is impossible to retrieve its information, within (eb) the life-time of a black hole, by externally coupling to (e&eb) the black hole qubits. However, authors show that, in principle, coupling to (e&eb) some of the internal degrees of freedom allows (eb) acquiring knowledge about the micro-state. Still, due to the trivial complexity of operations that can be performed, (eb) there is no time advantage over the collection of Hawking radiation and subsequent decoding. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Report number: LMU-ASC 03/16, MPP-2016-1 Cite as: arXiv:1605.01407 [hep-th] (or arXiv:1605.01407v1 [hep-th] for this version) **Universality of Black Hole Quantum Computing Gia Dvali, Cesar Gomez, Dieter Lust, Yasser Omar, and Benedikt Richter (Submitted on 4 May 2016) STOP** In this addendum to [arXiv: 1402.5674] two points are discussed. In the first additional evidence is provided (eb) for a dual connection between the geometric length of an Einstein-Rosen bridge and (e&eb) the computational complexity of the quantum state of the dual CFT's. The relation between growth of complexity and (e&eb) Page's “Extreme Cosmic Censorship” principle is also remarked on. The second point involves a gedanken experiment in which Alice measures (eb) a complete set of commuting observables at her end of (e) an Einstein-Rosen bridge is discussed. An apparent paradox is resolved by appealing to (e) the properties of GHZ tripartite entanglement. **Addendum to computational complexity and black hole horizons Leonard Susskind* Version of Record online: 18 JAN 2016 DOI: 10.1002/prop.201500093 STOP** Analyzing some well established facts, authors give a model-independent

parameterization of black hole quantum computing in terms of (e&*eb) a set of macro and micro quantities and their relations. These include the relations between the extraordinarily-small energy gap of (e) black hole qubits and important time-scales of (e) information-processing, such as, scrambling time and Page's time. Authors then show, confirming and extending previous results, that other systems of nature with (e&eb) identical quantum informatics features are (=) attractive Bose-Einstein systems at (eb) the critical point of quantum phase transition. Here authors establish a complete isomorphy between (e&eb) the quantum computational properties of these two systems. In particular, they show that the quantum hair of a critical condensate is (=) strikingly similar to (=) the quantum hair of a black hole. Irrespectively whether one takes the similarity between the two systems as a remarkable coincidence or as (=) a sign of a deeper underlying connection, the following is evident. Black holes are not unique in their way of (e) quantum information processing and they can manufacture black hole based quantum computers in labs by (e) taking advantage of quantum criticality. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Atomic Physics (physics. Atom-ph); Quantum Physics (quant-ph) Cite as: arXiv: 1601.01329 [hep-th] (or arXiv: 1601.01329v1 [hep-th] for this version) **Black Hole Based Quantum Computing in Labs and in the Sky Gia Dvali, Mischa Panchenko (Submitted on 6 Jan 2016) STOP Gia Dvali, Alexander Gußmann** show that the existence of black holes with (e&eb) classical skyrmion hair invalidates (e) standard proofs that global charges, such as (=) the baryon number, cannot be conserved by (e) a black hole. By carefully analyzing the standard arguments based on (e) a Gedankenexperiment in which a black hole is seemingly-unable to (e) return the baryon number that it swallowed, authors identify inconsistencies in this reasoning, which does not (e) take into the account neither the existence of skyrmion black holes nor the baryon/skyrmion correspondence. Authors then perform a refined Gedankenexperiment by incorporating (e) the new knowledge and show that (eb) no contradiction with conservation of baryon number takes place at (eb) any stage of black hole evolution. Analysis also indicates (eb) no conflict

between semi-classical black holes and (e&eb) the existence of baryonic gauge interaction arbitrarily-weaker than gravity. Next, authors study classical cross sections of a minimally-coupled massless probe scalar field scattered by (e) _ a skyrmion black hole. Authors investigate how the skyrmion hair manifests itself by (e) comparing this cross section with (e&eb) the analogous cross section caused by (e) a Schwarzschild black hole which has (e) the same ADM mass as (=) the skyrmion black hole. Here authors find an order-one difference in the positions of (e) the characteristic peaks in (eb) the cross sections. The peaks are shifted to smaller scattering angles when (e) the skyrmion hair is present. This comes from the fact that the skyrmion hair changes (e&eb) the near horizon geometry of the black hole when compared to (e&eb) a Schwarzschild black hole with same ADM mass. Authors keep the study of this second aspect general so that (e) the qualitative results obtained can also be applied to(e&eb) black holes with classical hair of different kind. Subjects: High Energy Physics - Theory (hep-th); High Energy Astrophysical Phenomena (astro-ph.HE); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU-ASC 18/16 Cite as: arXiv:1605.00543 [hep-th] (or arXiv:1605.00543v2 [hep-th] for this version) A preliminary discussion is given of (e) the prospects that gravitational-wave observations of (e) binary inspiral of black holes could reveal or constrain (e&eb) quantum modifications to (e) black hole dynamics, such as are required to (e) preserve postulates of quantum mechanics. Different proposals for such modifications are characterized by (e) different scales, and the size of these scales relative to (e&eb) those probed by observation of (e) inspiral signals is important in determining (eb) the feasibility of finding experimental signatures. Certain scenarios with strong quantum modifications in (eb) a region extending well outside the horizon are expected to modify (e&eb) classical evolution, and distort (e&eb) the near-peak gravitational wave signal, suggesting (eb) a search for departures from (e) waveforms predicted by general relativity. The near agreement of the GW150914 signals with (e&eb) such waveforms is discussed, and indicates (eb) constraints on some such scenarios. Important strategies for more precise future tests are 1) to develop more precise

predictions from (e) scenarios proposing quantum modifications, and 2) searching for observed deviations from (e) numerical relativity predictions via (e&eb) analysis of gravity wave data, particularly focussing on (e&eb) the signal region corresponding to plunge and merger. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1602.03622 [gr-qc] (or arXiv: 1602.03622v2 [gr-qc] for this version) **Gravitational wave tests of quantum modifications to black hole structure -- with post-GW150914 update Steven B. Giddings STOP** Direct LHC bounds on (e&eb) colored SUSY particles now corner naturalness more than the measured value of the Higgs mass does. Bounds on the gluino are of (e) particular importance, since it radiatively "sucks" up (e) the stop and Higgs-up soft masses. As a result, even models that easily accommodate a 125 GeV Higgs are almost as tuned as (=) the simplest version of SUSY, the MSSM: at best at the percent level. In this paper, authors further examine how current LHC results constrain (e) naturalness in three classes of models that may relax LHC bounds on (e&eb) sparticles: split families (e&eb), baryonic RPV (e&eb), and Dirac gauginos. In models of split families and bRPV, the bounds on the gluino are (=) only slightly reduced, resulting in (eb) a few percent tuning. In particular, having a natural spectrum in bRPV models typically implies (eb) that tops, Ws, and Zs are easily produced (eb) in the cascade decays of (e) squarks and gluinos. The resulting leptons and missing energy push (eb) the gluino mass limit above 1 TeV. Even when the gluino has a Dirac mass and does not contribute to the stop mass at one loop, tuning reappears in calculable models because there is no symmetry imposing the supersoft limit. Authors conclude that, even if sparticles are found at LHC-14, naturalness will not emerge triumphant. Subjects: High Energy Physics - Phenomenology (hep-ph)DOI:10.1007/JHEP03(2014)022 Cite as:rXiv:1309.3568 [hep-ph] (or arXiv:1309.3568v2 [hep-ph] for this version) **The Last Vestiges of Naturalness Asimina Arvanitaki, Masha Baryakhtar, Xinlu Huang, Ken Van Tilburg, Giovanni Villadoro STOP** The lack of evidence for new physics beyond (e) the standard model at the LHC points to (e) a paucity of new particles near the weak scale. This suggests that the weak scale is tuned and that supersymmetry, if present at

all, is realized (eb) at higher energies. The measured Higgs mass constrains (e) the scalar sparticles to be below 10^5 TeV, while (e) gauge coupling unification favors (eb) Higgsinos below 100 TeV. Nevertheless, in many models gaugino masses are (=) suppressed and remain within reach of the LHC. Tuning the weak scale and (e&eb) the renormalization group evolution of (e) the scalar masses constrain (e) Split model building. Due to the small gaugino masses, either the squarks or the up-higgs often run tachyonic; in the latter case, successful electroweak breaking requires heavy higgsinos near the scalar sparticles. We discuss the consequences of tuning the weak scale and the phenomenology of several models of Split supersymmetry including anomaly mediation, U(1)_(B-L) mediation, and Split gauge mediation. Subjects: High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Experiment (hep-ex) DOI: 10.1007/JHEP02 (2013)126 Cite as: arXiv: 1210.0555 [hep-ph] (or arXiv: 1210.0555v2 [hep-ph] for this version) **Mini-Split Asimina Arvanitaki, Nathaniel Craig, Savas Dimopoulos, Giovanni Villadoro STOP** Light states associated with (e&eb) the hierarchy problem affect (e&eb) the Higgs LHC production and decays. Authors illustrate this within the MSSM and two simple extensions applying (e&eb) the latest bounds from (e) LHC Higgs searches. Large deviations in the Higgs properties are expected in (eb) a natural SUSY spectrum. The discovery of a non-Standard-Model Higgs may signal (eb) the presence of light stops accessible at (eb) the LHC. Conversely, the more the Higgs is Standard-Model-like, (eb) the more tuned the theory becomes. Taking the ratio of different Higgs decay channels at the LHC cancels (er&eb) the leading (eb) QCD uncertainties and potentially improves (eb) the accuracy in Higgs coupling measurements to (e) the percent level. This may lead to (eb) the possibility of doing precision Higgs physics at (eb) the LHC. Finally, authors entertain the possibility that the ATLAS excess around (e&eb) 125 GeV persists with (e&eb) a Higgs production cross-section that is enhanced compared to (e&eb) the SM. This increase can only be accommodated in extensions of (e) the MSSM and it may suggest (eb) that stops lie below 400 GeV, likely within (eb) reach of next year's LHC run. Subjects: High Energy Physics - Phenomenology (hep-ph) DOI: 10.1007/JHEP02 (2012)144 Cite

as: arXiv: 1112.4835 [hep-ph] (or arXiv: 1112.4835v2 [hep-ph] for this version) **A Non Standard Model Higgs at the LHC as a Sign of Naturalness Asimina Arvanitaki, Giovanni Villadoro STOP** It has recently been suggested that the presence of a plenitude of light axions, an Axiverse, is (=) evidence for the extra dimensions of string theory. Authors discuss the observational consequences of (e) these axions on astrophysical black holes through (e&eb) the Penrose superradiance process. When an axion Compton wavelength is comparable to (e&eb) the size of a black hole, the axion binds to (e&eb) the black hole "nucleus" forming (eb) a gravitational atom in the sky. The occupation number of superradiant atomic levels fed by (e) the energy and angular momentum of the black hole grows exponentially. The black hole spins down and an axion Bose-Einstein condensate cloud forms (eb) around it. When the attractive axion self-interactions become (=) stronger than the gravitational binding energy, (eb) the axion cloud collapses, a phenomenon known in condensed matter physics as "Bosenova". The existence of axions is first diagnosed by (e) gaps in the mass vs (e&eb) spin plot of astrophysical black holes. For young black holes the allowed values of (e) spin are quantized, giving rise to (eb) "Regge trajectories" inside the gap region. The axion cloud can also be observed directly either through (e&eb) precision mapping of the near horizon geometry or through (e&eb) gravitational waves coming from (e) the Bosenova explosion, as well as axion transitions and (e&eb) annihilations in (eb) the gravitational atom. Estimates suggest (eb) that these signals are detectable in (eb) upcoming experiments, such as Advanced LIGO, AGIS, and LISA. Current black hole spin measurements imply (eb) an upper bound on (e&eb) the QCD axion decay constant of (e) 2 x 10^17 GeV, while (e) Advanced LIGO can detect (e&eb) signals from (e) a QCD axion cloud with (e&eb) a decay constant as low as (=) the GUT scale. Authors finally discuss the possibility of observing the gamma-rays associated with the Bosenova explosion and, perhaps, the radio waves from axion-to-photon conversion for the QCD axion. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Journal

reference:Phys.Rev.D83:044026, 2011 DOI: 10.1103/PhysRevD.83.044026 Cite as: arXiv: 1004.3558 [hep-th] (or arXiv: 1004.3558v2 [hep-th] for this version) **Exploring the String Axiverse with Precision Black Hole Physics Asimina Arvanitaki, Sergei Dubovsky (Submitted on 20 Apr 2010 (v1), last revised 16 Aug 2011 (this version, v2)) STOP Skyrmion Black Hole Hair: Conservation of Baryon Number by Black Holes and Observable Manifestations Gia Dvali, Alexander Gußmann STOP Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust** discuss BMS supertranslations both at (eb) null-infinity and on the horizon for the case of the Schwarzschild black hole. Authors show that both kinds of supertranslations lead to (eb) infinitely many gapless physical excitations. On this basis authors construct a quotient algebra using (e) suited superpositions of (e) both kinds of transformations which cannot be compensated by (e) an ordinary BMS-supertranslation and therefore are (=) intrinsically due to (e) the presence of an event horizon. Authors show that these quotient transformations are (=) physical and generate (eb) gapless excitations on the horizon that can account for (e) the gravitational hair as well as for (e) the black hole entropy. They identify the physics of these modes as associated with (e&eb) Bogolioubov-Goldstone modes due to quantum criticality. Classically the number of these gapless modes is (=) infinite. However, authors show that due to (e) quantum criticality the actual amount of information-carriers becomes (=) finite and consistent with Bekenstein entropy. Although they only consider the case of Schwarzschild geometry, the arguments are extendable to(e&eb) arbitrary space-times containing event horizons. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) Report number: LMU--ASC 05/16, MPP--2016—3 Cite as: arXiv: 1601.03725 [hep-th] (or arXiv: 1601.03725v1 [hep-th] for this version) **Gravitational Black Hole Hair from Event Horizon Supertranslations Artem Averin, Gia Dvali, Cesar Gomez, Dieter Lust STOP Gia Dvali, Cesar Gomez, Nico Wintergerst** suggest that holography can be formulated in terms of (e&eb) the information capacity of (e) the St\"uckelberg degrees of freedom that maintain

gauge invariance of (e) the theory in the presence of (e) an information boundary. These St\"uckelbergs act as (e&eb) qubits that account for a certain fraction of (e) quantum information. Their information capacity is measured by (e) the ratio of the inverse St\"uckelberg energy gap to (e) the size of the system. Systems with (e&eb) the smallest gap are (=) maximally holographic. For massless gauge systems this information measure is (=) universally equal to (=) the inverse coupling evaluated at the systems' length scale. In this language it becomes very transparent why the St\"uckelberg information capacity of black holes saturates (e full) the Bekenstein bound and accounts for (e) the entire information of the system. The physical reason is that the strength of quantum interaction is bounded from below by the gravitational coupling, which scales as (=) area. Observing the striking similarity between the scalings of the energy gap of the boundary St\"uckelberg modes and (e&eb): (AC=DC=0) the Bogoliubov modes of critical many-body systems, authors establish a connection between holography and (e&eb) quantum criticality through (e&eb) the correspondence between these modes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1511.03525 [hep-th] (or arXiv: 1511.03525v1 [hep-th] for this version) **Stückelberg Formulation of Holography Gia Dvali, Cesar Gomez, Nico Wintergerst STOP B**lack hole entropy, denoted by N, in (semi)classical limit is (=) infinite. This scaling reveals (eb) a very important information about the qubit degrees of (e) freedom that carry black hole entropy. Namely, the multiplicity of qubits scales as N, whereas (e) their energy gap and their coupling as 1/N. Such a behavior is indeed exhibited by (e) Bogoliubov-Goldstone degrees of freedom of (e) a quantum-critical state of N soft gravitons (a condensate or a coherent state) describing (eb) the black hole quantum portrait. They can be viewed as (=) the Goldstone modes of (e) a broken symmetry acting on (e&eb) the graviton condensate. In this picture Minkowski space naturally emerges as (=) a coherent state of infinite-N gravitons of (e) infinite wavelength and it carries (e&eb) an infinite entropy. In this paper authors **Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015)** ask what is the geometric meaning (if any) of (e) the classical limit of this symmetry. Authors argue that the infinite-

N limit of Bogoliubov-Goldstone modes of critical graviton condensate is described by (e) recently-discussed classical BMS super-translations broken by the black hole geometry. However, the full black hole information can only be recovered for (e) finite N, since the recovery time becomes (=) infinite in classical limit in which N is (=) infinite. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph) DOI: 10.1016/j.physletb.2015.11.073 Report number: MPP-2015-215, LMU-ASC 59/15 Cite as: arXiv: 1509.02114 [hep-th] (or arXiv: 1509.02114v1 [hep-th] for this version) **Classical Limit of Black Hole Quantum N-Portrait and BMS Symmetry Gia Dvali, Cesar Gomez, Dieter Lüst (Submitted on 7 Sep 2015) STOP** Recent ideas about understanding physics of black hole information-processing in terms of (e&eb) quantum criticality allow us to (eb) implement black hole mechanisms of quantum computing within (eb) critical Bose-Einstein systems. The generic feature, uncovered both by analytic and numeric studies, is (=) the emergence at the critical point of (e) gapless weakly-interacting modes, which act as (=) qubits for information-storage at a very low energy cost. These modes can be effectively described in terms of (e&eb) either Bogoliubov or Goldstone degrees of freedom. The ground-state at the critical point is (=) maximally entangled and far from being classical. We confirm this near-critical behavior by a new analytic method. **Gia Dvali, Mischa Panchenko** compute growth of entanglement and show its consistency with (e&eb) black hole type behavior. On the other hand, in the over-critical regime the system develops (eb) a Lyapunov exponent and scrambles quantum information very fast. By, manipulating (e&eb) the system parameters externally, authors can put it in and out of various regimes and in this way control (e&eb) the sequence of information storage and processing. By using gapless Bogoliubov modes as (=) control qubits they design some simple logic gates. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Cite as: arXiv: 1507.08952 [hep-th] (or arXiv: 1507.08952v1 [hep-th] for this version) **Black Hole Type Quantum**

**Computing in Critical Bose-Einstein Systems Gia Dvali, Mischa Panchenko (Submitted on 31 Jul 2015) STOP Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst** establish a fundamental connection between quantum criticality of a many-body system, such as Bose-Einstein condensates, and (e&eb) its capacity of information-storage and processing. For deriving the effective theory of modes in the vicinity of the quantum critical point they develop a new method by mapping a Bose-Einstein condensate of N-particles onto (e&eb) a sigma model with a continuous global (pseudo)symmetry that mixes (e&eb) bosons of different momenta. The Bogolyubov modes of the condensate are mapped onto (e&eb) the Goldstone modes of the sigma model, which become (=) gapless at the critical point. These gapless Goldstone modes are (=) the quantum carriers of information and entropy. Analyzing their effective theory, authors observe the information-processing properties strikingly similar to (=) the ones predicted by the black hole portrait. The energy cost per qubit of information-storage vanishes in (eb) the large-N limit and the total information-storage capacity increases with (e&eb) N either exponentially or as a power law. The longevity of information-storage also increases with (e&eb) N, whereas (e) the scrambling time in (eb) the over-critical regime is controlled by (e&eb) the Lyapunov exponent and scales logarithmically with (e&eb) N. This connection reveals (eb) that the origin of black hole information storage lies in (eb) the quantum criticality of the graviton Bose-gas, and that much simpler systems that can be manufactured in table-top experiments can exhibit (eb) very similar information-processing dynamics. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 92, 125002 (2015) DOI: 10.1103/PhysRevD.92.125002 Cite as: arXiv:1507.02948 [hep-th] (or arXiv:1507.02948v1 [hep-th] for this version) **Nambu-Goldstone Effective Theory of Information at Quantum Criticality Gia Dvali, Andre Franca, Cesar Gomez, Nico Wintergerst STOP Gia Dvali, Cesar Gomez** show that a positive cosmological constant is (=) incompatible with the quantum-corpuscular resolution of de Sitter metric in form of (e) a coherent state. The reason is very general and is due to (e) the

quantum self-destruction of (e) the coherent state because of (e) the scattering of constituent graviton quanta. This process creates an irreversible quantum clock, which precludes (e) eternal de Sitter. It also eliminates (e) the possibility of Boltzmann brains and Poincare recurrences. This effect is expected to be (=) part of any microscopic theory that takes into account (e) the quantum corpuscular structure of the cosmological background. This observation puts the cosmological constant problem in a very different light, promoting it, from a naturalness problem, into (e&eb) a question of quantum consistency. We are learning that quantum gravity cannot tolerate (e) exceedingly-classical sources. Subjects: High Energy Physics - Theory (hep-th); Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1412.8077 [hep-th]

(or arXiv:1412.8077v1 [hep-th] for this version) **Quantum Exclusion of Positive Cosmological Constant? Gia Dvali, Cesar Gomez (Submitted on 27 Dec 2014) STOP G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger** establish a connection between the ultra-Planckian scattering amplitudes in field and string theory and (e&eb) unitarization by black hole formation in these scattering processes. **Using as a guideline an explicit microscopic theory in which the black hole represents (eb) a bound-state of many soft gravitons at (eb) the quantum critical point, they were able to identify and compute (e&eb) a set of perturbative amplitudes relevant for (e) black hole formation. These are the tree-level N-graviton scattering S-matrix elements in (eb) a kinematical regime (called classicalization limit) where (e) the two incoming ultra-Planckian gravitons produce (eb) a large number N of soft gravitons. Authors compute these amplitudes by using (e) the Kawai-Lewellen-Tye relations, as well as scattering equations and (e&eb) string theory techniques. Authors discover that this limit reveals (eb) the key features of the microscopic corpuscular black hole N-portrait. In particular, the perturbative suppression factor of a N-graviton final state, derived from (e) the amplitude, matches the non-perturbative black hole entropy when (e) N reaches the quantum criticality value, whereas final states with different value of N are (=) either suppressed or excluded by non-perturbative**

**corpuscular physics. Thus they identify the microscopic reason behind the black hole dominance over other final states including (e) non-black hole classical object. In the parameterization of the classicalization limit the scattering equations can be solved exactly allowing (eb) us to obtain closed expressions for the high-energy limit of the open and closed superstring tree-level scattering amplitudes for a generic number N of external legs. Authors demonstrate matching and complementarity between the string theory and (e&eb) field theory** in different large-s and large-N regimes. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1016/j.nuclphysb.2015.02.004 Report number: MPP-2014-320, LMU-ASC 52/14 Cite as: arXiv:1409.7405 [hep-th] (or arXiv:1409.7405v2 [hep-th] for this version) **Black Hole Formation and Classicalization in Ultra-Planckian 2 -> N Scattering G. Dvali, C. Gomez, R.S. Isermann, D. Lust, S. Stieberger STOP** Gravitational backgrounds, such as black holes, AdS, de Sitter and inflationary universes, should be viewed as (=) composite of N soft constituent gravitons. It then follows that such systems are close to (e) quantum criticality of graviton Bose-gas to Bose-liquid transition. Generic properties of the ordinary metric description, including (e) geodesic motion or particle-creation in the background metric, emerge as (=) the large-N limit of quantum scattering of (e) constituent longitudinal gravitons. **Gia Dvali, Cesar Gomez** show that this picture correctly accounts for (e) physics of large and small black holes in AdS, as well as reproduces (eb) well-known inflationary predictions for cosmological parameters. However, it anticipates new effects not captured by (e) the standard semi-classical treatment. In particular, **Gia Dvali, Cesar Gomez** predict observable corrections that are sensitive to (e&eb) the inflationary history way beyond last 60 e-foldings. They derive an absolute upper bound on the number of e-foldings, beyond (e) which neither de Sitter nor inflationary Universe can be approximated by (e) a semi-classical metric. However, they could in principle persist in a new type of (e) it quantum eternity state. Authors discuss implications of this phenomenon for (e&eb) the cosmological constant problem. Subjects: High Energy Physics - Theory (hep-th); Cosmology and

Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Quantum Physics (quant-ph) DOI: 10.1088/1475-7516/2014/01/023 Cite as: arXiv: 1312.4795 [hep-th] (or arXiv:1312.4795v1 [hep-th] for this version) **Quantum Compositeness of Gravity: Black Holes, AdS and Inflation Gia Dvali, Cesar Gomez STOP** As quantum information science approaches the goal of constructing quantum computers, understanding loss of information through (e&eb) decoherence becomes (=) increasingly important. The information about a system that can be obtained from (e) its environment can facilitate quantum control and error correction. Moreover, observers gain (eb) most of their information indirectly, by (e) monitoring (primarily photon) environments of (e) the "objects of interest." Exactly how this information is inscribed in the environment is essential for (e) the emergence of "the classical" from (e) the quantum substrate. In this paper, **Robin Blume-Kohout, Wojciech H. Zurek** examine how many-qubit (or many-spin) environments can store (e) information about a single system. The information lost to (e) the environment can be stored redundantly, or it can be encoded in (eb) entangled modes of the environment. Authors go on to show that randomly chosen states of the environment almost always encode (e) the information so that (e) an observer must capture (e) a majority of the environment to deduce (eb) the system's state. Conversely, in the states produced by (e) a typical decoherence process, information about a particular observable of the system is stored redundantly. This selective proliferation of "the fittest information" (known as Quantum Darwinism) plays a key role in (eb) choosing the preferred, effectively classical observables of (e) macroscopic systems. The developing appreciation that the environment functions not just as (=) a garbage dump, but as a communication channel, is extending understanding of the environment's role in (e&eb) the quantum-classical transition beyond (e) the traditional paradigm of decoherence. Keywords Decoherence quantum Darwinism quantum information quantum-classical transition **Foundations of Physics November 2005, Volume 35, Issue 11, pp 1857-1876 A Simple Example of "Quantum Darwinism": Redundant Information Storage in Many-Spin Environments Robin Blume-Kohout, Wojciech H. Zurek STOP Harold Ollivier, David**

**Poulin, and Wojciech H. Zurek** study the emergence of objective properties in open quantum systems. In our analysis, the environment is promoted from a passive role of a reservoir selectively destroying quantum coherence to an active role of amplifier selectively **proliferating information about the system. We show that only preferred pointer states of the system can leave a redundant and therefore easily detectable imprint on the environment. Observers who—as is almost always the case—discover the state of the system indirectly (by probing a fraction of its environment) will find out only about the corresponding pointer observable. Many observers can act in this fashion independently and without perturbing the system. They will agree about its state. In this operational sense, preferred pointer states exist objectively. Received 24 July 2003 DOI:http://dx.doi.org/10.1103/PhysRevLett.93.220401 Objective Properties from Subjective Quantum States: Environment as a Witness Harold Ollivier, David Poulin, and Wojciech H. Zurek Phys. Rev. Lett. 93, 220401 – Published 22 November 2004 STOP** Quantum Darwinism recognizes that we—the observers—acquire our information about the "systems of interest" indirectly from their imprints on the environment. **Michael Zwolak, H. T. Quan, and Wojciech H. Zurek** show that information about a system can be acquired from a mixed-state, or hazy, environment, but the storage capacity of an environment fragment is suppressed by its initial entropy. In the case of good decoherence, the mutual information between the system and the fragment is given solely by the fragment's entropy increase. For fairly mixed environments, this means a reduction by a factor 1−h, where h is the haziness of the environment, i.e., the initial entropy of an environment qubit. Thus, even such hazy environments eventually reveal the state of the system, although now the intercepted environment fragment must be larger by ~ (1−h)−1 to gain the same information about the system. DOI:http://dx.doi.org/10.1103/PhysRevLett.103.110402 **Quantum Darwinism in a Mixed Environment Michael Zwolak, H. T. Quan, and Wojciech H. Zurek Phys. Rev. Lett. 103, 110402 – Published 8 September 2009** We study quantum Darwinism—the redundant recording of information about the preferred states of a decohering system by its environment—for an object illuminated by a

blackbody. In the cases of point-source and isotropic illumination, we calculate the quantum mutual information between the object and its photon environment. We demonstrate that this realistic model exhibits fast and extensive proliferation of information about the object into the environment and results in redundancies orders of magnitude larger than the exactly soluble models considered to date. DOI:http://dx.doi.org/10.1103/PhysRevLett.105.020404 **Quantum Darwinism in an Everyday Environment: Huge Redundancy in Scattered Photons C. Jess Riedel and Wojciech H. Zurek Phys. Rev. Lett. 105, 020404 – Published 9 July 2010 STOP** Measurements transfer information about a system to the apparatus and then, further on, to observers and (often inadvertently) to the environment. I show that even imperfect copying essential in such situations restricts possible unperturbed outcomes to an orthogonal subset of all possible states of the system, thus breaking the unitary symmetry of its Hilbert space implied by the quantum superposition principle. Preferred outcome states emerge as a result. They provide a framework for "wave-packet collapse," designating terminal points of quantum jumps and defining the measured observable by specifying its eigenstates. In quantum Darwinism, they are the progenitors of multiple copies spread throughout the environment—the fittest quantum states that not only survive decoherence, but subvert the environment into carrying information about them—into becoming a witness. DOI:http://dx.doi.org/10.1103/PhysRevA.76.052110 **Quantum origin of quantum jumps: Breaking of unitary symmetry induced by information transfer in the transition from quantum to classical Wojciech Hubert Zurek Phys. Rev. A 76, 052110 – Published 16 November 2007 STOP** The Born rule, a foundational axiom used to deduce probabilities of events from **wavefunctions, is indispensable in the everyday practice of quantum physics. It is also key in the quest to reconcile the ostensibly inconsistent laws of the quantum and classical realms, as it confers physical significance to reduced density matrices, the essential tools of decoherence theory. Following Bohr's Copenhagen interpretation, textbooks postulate the Born rule outright. However, recent attempts to derive it from (e) other quantum principles have been successful, holding promise for (e)**

**simplifying and clarifying the quantum foundational bedrock. A major family of derivations is based on envariance, a recently discovered symmetry of (e) entangled quantum states. Here, authors identify and experimentally test three premises central to (e) these envariance-based derivations, thus demonstrating, (eb) in the microworld, the symmetries from (e) which the Born rule is derived. Further, they demonstrate envariance in a purely local quantum system, showing (eb) its independence from relativistic causality.Subjects: Quantum Physics (quant-ph)Cite as:arXiv:1604.01471 [quant-ph] (or arXiv:1604.01471v1 [quant-ph] for this version) Quantum probabilities from quantum entanglement: Experimentally unpacking the Born rule Jérémie Harris, Frédéric Bouchard, Enrico Santamato, Wojciech H. Zurek, Robert W. Boyd, Ebrahim Karimi**

Integrating $[m_4]_4(t, a_4)$ with respect to a from 0 to ∞, we have

$$\frac{dM_4(t)}{dt} = \int_0^\infty [mt]_4(t, a_4)da_4$$

$$= -\int_0^\infty [m_a]_4\,(t, a_4)da_4 - \int_0^\infty \mu(a_4)m(t, a_4)da_4$$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f^{'}{}_{4*}]_4 \int_0^\infty p_4^*(a_4)da_4\right)M_4 - [f_{2*}^4]_4 \int_0^\infty p_4^*(a_4)da_4 W_4$$

$= m_4(t, 0) - \int_0^\infty \mu_4(a_4)\, m_4(t, a_4)da_4$

$$-\left([f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}^{i}]_4 P_4^*\right)M_4 - [f_{2*}^{i}]_4 P_4^* W_{\cdot 4}$$

From (30), we arrive at the equations

$$\frac{dW_4(t)}{dt} = \int_0^{\infty} ([h_*]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4$$

$$-\left([f_{4*}] + [f_{2*}]_4 + \left[f_{4*}^{'}\right]_4 P_4^*\right) M_4 + \left(\frac{\left[h_*^{'}\right]_4}{[h_*]^4}\, P_4^*(0) - f_{2*}^{'} P_4^*\right) W_4,$$

(34)

$$\frac{dW_4(t)}{dt}\left[g_{*}^{'}\right]_4 Q_4^*\, M_{.4}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{r*}]_4\beta_4(a_4) - \mu_4(a_4))m_4(t,a_4)da_4 = ([h_*]_4\beta_4(\theta_4) - \mu_4(\theta_4))M_{,4}$$

Where $0< \theta_4 = \theta_4(t) < \infty$.

And if $a_4^2\,[a_4^2]_4<4\,[a_2]_4<1\,a_2 e$,

$\mathrm{Re}\,\lambda_4\left[\lambda_4\right]_4 = \mathrm{Re}\lambda_2\left[\lambda_2\right]_4 = -\frac{4}{2}a_4[a2\ ]_4 < 0.$

Let $\Delta=\max_{2<u<\infty}\left|h_*\,[h_*]_4\beta_4(a_4) - \mu_4(a_4)\right|$. Then$|n_4(t)|$ ⊠⊿∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_4^2 <4\, a_2 e$ and $[a_4] > 2\Delta$,or $[a_4^2]$ ⊠⊿$[a_2]e$ and $(a_4 - \sqrt{(a_4^2)_4 - 4(a_2)_4 e) > 2\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+4}\right)$ and $\beta_4(a_4) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty} \beta_4(a_4) e^{-\int_0^a \mu_4(\tau_4) + d\tau_4 da} = \frac{4}{(m + \mu_0 + 1)^2} + \frac{2}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_4 + [\mu_0]$ 0.5. Assume$[f_4](x_4)= K_4 x_4$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_4(x_4)= [0x_4]/(1 + x_4)$; the functional responses $[[f_2](x) = [0]\, x_4/(1 + x_4)$and $h_4(x_4)=(1+\omega_4 x_4)/(1 + x_4)$(See Holling [7] and Rashevsky [13])

Equation (44),now has the form of

$$\frac{m_4 + [\mu_0] + [K][P_0])}{[m_4]_4 + (\mu_0)_4 + [K][P_0])\frac{3}{4}\,[K]_4[P_0]\frac{3}{4}} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(4 + P_0) < b_4$ i e $[P_0] < b/(\delta_4 - b_4)$then the equilibrium$(P_0)$ $[P_0]_4(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_4$

$[f_{4*}] = K_4\, P^*\, [f_{2*}] = \frac{\sigma Q_4^*}{4+Q_4^*} \quad [h_*] = \frac{4+\omega_4 Q_4^*}{4+Q_4^*}$

$[f_{4*}] = K_4\, [f_{2*}^4] = \frac{\sigma_4}{(4+Q_4^*)^2}\, [g_*^4] = \frac{\sigma_4}{(4+P_4^*)^2} \quad [h_*^4] = \frac{\omega_4 - 4}{(4+Q_4^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0$, K, b, $\delta$, and $\omega$ need to satisfy

$$\frac{\omega_4}{(m + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma)^2} + \frac{2\omega_4}{([m_4]_4 + [\mu_0]_4 + 1 + \frac{K_4 b_4}{\delta_4 + b_4} + \sigma_4)^3} < 1$$

And

$$\frac{1}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^2}+\frac{2\omega_4}{(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_4+[\mu_0]_4+1+\frac{K_4b_4}{\delta_4-b_4}+(m_4+[\mu_0]_4+\frac{K_4b_4}{\delta_4-b_4}+\sigma_4)\,Q_4^*)^3$$

$$=m_4+[\mu_0]_4+3+\frac{K_4b_4}{\delta_4-b_4}(4+Q^*)(1+\omega_4Q_4^*$$

Now

$$\left[a_4^2\right]-1\left[a_2\right]e=\left(\frac{2K_4b_4}{\delta_4-b_4}\right)^2+\frac{1K_4b_4\sigma_4\,Q_4^*}{(\delta_4-b_4)(1+Q_4^*)}+\frac{\sigma_4^2Q_4^{*2}}{(1+Q_4^*)^2}$$

$$\frac{1\,b_4\sigma_4\,(\delta_4-b)Q^*}{\delta_4(1+Q_4^*)^2}-\frac{1\,(\delta_4-b_4)^2(1-\omega_4)Q_4^*}{\delta_4(1+\omega_4Q_4^*)(1+Q_4^*)}P^*(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_4\,ae_4^{mu}\;[\mu_0]_4-\frac{a_4}{a_4+4}$$

Hence if $[a_4^2>4(a_2)_4e$ and

$$[a_4] = \frac{2K_4 b_4}{\delta_4 - b_4} + \frac{\sigma_4 \, Q_4^*}{1 + Q_4^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_4]_4 = [f_{4*}]_4 + [f_{2*}]_4 + [f_{4*}] P_4^*$$

$$[a_2]_4 = \left[f_{2*}^4\right] P_4^* - \frac{h_*^4}{h_*} P^*(0)$$

$$e = [g_*^4]_4 \, Q_4^*$$

$$\eta_4(t) = \left[h_*\right]_4 \beta_4 \left(\theta_4(t)\right) - \mu_4(\theta(t)),$$

The system(34)can be written as the following system:

$$[M(t)]\left[a_4\right] M_4 - [a_2]_4 W_4 - [n(t)] M_4$$

$$w_4(t) = eM_4$$

Write

$$A = \begin{pmatrix} -[a_4]_4 - [a_2]_4 & \\ e & 0 \end{pmatrix}$$

$$B = (t) = \begin{pmatrix} [\eta(t)]_4 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_4 \\ W_4 \end{pmatrix}$$

The solution of (32)is

$$X_4(t) = e^{A,t}\,[X]_4(0) + \int_0 e^{At}$$

**Some of the other conceptual and thematic and discursive factorial representationalities to which the model is applicable are:every one-dimensional quantum mechanical Hamiltonian H1 can have (e) a partner H2 such that H1 and H2 taken together may be viewed as (=) the components of a supersymmetric Hamiltonian. The term 'supersymmetric Hamiltonian' is taken to mean (eb) a Hamiltonian defined in terms of (e&eb)charges that obey (eb)the same algebra as that of (e) the generators of supersymmetry in field theory. The consequences of this symmetry for (er) the spectra of H1 and H2 are explored. It is shown how the supersymmetric pairing may be utilised to eliminate (e) the ground state of H1, or add a state below the ground state of H1 or maintain the spectrum of H1. It is also explicitly demonstrated that the supersymmetric pairing may be used to generate (eb) a class of anharmonic potentials with (e&eb) exactly specified spectra. The complete spectrum of an anharmonic potential so generated consists of (e) all the eigenstates of the simple harmonic oscillator and, in addition, a ground state at a specified energy E which lies arbitrarily below (eb) the E=1/2 ground state of the harmonic oscillator. Supersymmetric quantum mechanics of one-dimensional systems C V Sukumar Journal of Physics A: Mathematical and General, Volume 18, Number 15 STOP Authors examine the properties and consequences of (e&eb) pseudo-supersymmetry for (e) quantum systems admitting (e) a pseudo-Hermitian Hamiltonian. Authors explore the Witten index of (e) pseudo-**

**supersymmetry and show that every pair of diagonalizable (not necessarily Hermitian) Hamiltonians with (e&eb) discrete spectra and (e&eb) real or complex-conjugate pairs of (e) eigenvalues are (=) isospectral and have (e) identical degeneracy structure except (e) perhaps for the zero eigenvalue if and only if (e) they are pseudo-supersymmetric partners. This implies (eb) that pseudo-supersymmetry is (=) the basic framework for (e) generating non-Hermitian PT-symmetric and (e&eb) non-PT-symmetric Hamiltonians with (e&eb) a real spectrum via (e&eb) a Darboux transformation, and shows (eb) that every diagonalizable Hamiltonian H with(e&eb) a discrete spectrum and (e&eb) real or complex-conjugate pairs of (e) eigenvalues may be factored as (=) H=L♯L where L is (=) a linear operator with pseudo-adjoint L♯. In particular, this factorization applies to (e&eb) PT-symmetric and Hermitian Hamiltonians. The non-degenerate two-level systems provide (eb) a class of Hamiltonians that are pseudo-Hermitian. Authors demonstrate the implications of general results for (e&eb) this class in some detail. Copyright © 2002 Elsevier Science B.V. All rights reserved. Nuclear Physics B Volume 640, Issue 3, 30 September 2002, Pages 419–434Pseudo-supersymmetric quantum mechanics and isospectral pseudo-Hermitian Hamiltonians Ali Mostafazadeh STOP Michele Caselle, Alessandro Nada, Marco Panero present a high-precision lattice calculation of (e) the equation of state in the confining phase of (e) SU(2) Yang-Mills theory. Authors show that the results are described very well by (e) a gas of massive, non-interacting glueballs, provided (e) one assumes an exponentially growing Hagedorn spectrum. The latter can be derived within (eb) an effective bosonic closed-string model, leading to (eb) a parameter-free theoretical prediction, which is in perfect agreement with (=) proposed lattice results. Furthermore, when applied to (e&eb) SU (3) Yang-Mills theory, this effective model accurately describes (eb) the lattice results reported by Bors\'anyi et al. in JHEP 07 (2012) 056. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 07 (2015) 143 DOI: 10.1007/JHEP07 (2015)143 Cite as: arXiv:**

**1505.01106 [hep-lat] (or arXiv: 1505.01106v2 [hep-lat] for this version) Hagedorn spectrum and thermodynamics of SU (2) and SU (3) Yang-Mills theories Michele Caselle, Alessandro Nada, Marco Panero STOP B. de Wit J. Hoppe, H. Nicolai study the quantum-mechanical properties of (e) a supermembrane and examine the nature of (e) its ground state. A supersymmetric gauge theory of (e) area-preserving transformations provides (eb) a convenient framework for this study. The supermembrane can be viewed as (=) a limiting case of a class of models in supersymmetric quantum mechanics. Its mass does not depend on (=) the zero modes and vanishes only if (e) the wave function is (=) a singlet under (e&eb) supersymmetry transformations of (e) the nonzero modes. Authors exhibit (eb) the complexity of (e) the supermembrane ground state and (e&eb) examine various truncations of these models. None of these truncations has massless states. Copyright © 1988 Published by Elsevier B.V. Nuclear Physics B Volume 305, Issue 4, 28 December 1988: Pages 545-581: On the quantum mechanics of supermembranes Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more links. B. de Wit J. Hoppe, H. Nicolai STOP In this note we explain how world-volume geometries of (e) D-branes can be reconstructed within (eb) the microscopic framework where (e) D-branes are described through (e&eb) boundary conformal field theory. Authors extract the (non-commutative) world-volume algebras from (e) the operator product expansions of (e) open string vertex operators. For branes in a flat background with (e&eb) constant non-vanishing B-field, the operator products are computed perturbatively to (e) all orders in the field strength. Resulting series coincides with (=) Kontsevich's presentation of the Moyal product. After extending these considerations to (e) fermionic fields authors conclude with some remarks on the generalization of approach to (e&eb) curved backgrounds. D-branes and deformation quantization Volker Schomerus1,2 Published 25 July 1999 Journal of High Energy Physics, Volume 1999, JHEP06(1999)STOP Authors employ the path integral approach developed in Gangopadhyay S. and Scholtz F. E., Phys. Rev.**

**Lett., 102 (2009) 241602 to discuss the (generalized) harmonic oscillator in (eb) a noncommutative plane. The action for (e) this system is derived in (e&eb) the coherent state basis with (e&eb) additional degrees of freedom. From this the action in the coherent state basis without (e) any additional degrees of freedom is obtained. This gives (eb) the ground-state spectrum of the system. Authors then employ the exact renormalization group approach to show (eb) that equivalence can be constructed between this (noncommutative) system and (e&eb) a commutative system Path integral action and exact renormalization group dualities for quantum systems in noncommutative plane Sunandan Gangopadhyay Published 25 June 2015 • Copyright © EPLA, 2015 STOP Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini present a lattice study of (e)the equation of state in Yang-Mills theory based on (e) the exceptional G(2) gauge group. As is well-known, at zero temperature this theory shares (e&eb) many qualitative features with real-world QCD, including (e) the absence of colored states in the spectrum and dynamical string breaking at (e) large distances. In agreement with previous works, authors show that at finite temperature this theory features (eb) a first-order deconfining phase transition, whose nature can be studied by (e) a semi-classical computation. Authors also show that the equilibrium thermodynamic observables in the deconfined phase bear striking quantitative similarities with (=) those found in SU (N) gauge theories: in particular, these quantities exhibit (eb) nearly perfect proportionality to the number of gluon degrees of freedom, and (e&eb) the trace anomaly reveals (eb) a characteristic quadratic dependence on (e&eb) the temperature, also observed in SU (N) Yang-Mills theories (both in four and in three spacetime dimensions). Authors compare our lattice data with (e&eb) analytical predictions from (e) effective models, and discuss their implications for (e) the deconfinement mechanism and (e&eb) high-temperature properties of (e) strongly interacting, non-supersymmetric gauge theories. Results give (eb) strong evidence for (e) the conjecture that the thermal deconfining transition is governed by (e&eb) a universal mechanism, common to (e) all simple gauge groups.**

**Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th); Nuclear Theory (nucl-th) Journal reference: JHEP 03 (2015) 057 DOI: 10.1007/JHEP03 (2015)057 Report number: DESY-14-146; IFT-UAM/CSIC-14-076 Cite as: arXiv: 1409.8305 [hep-lat] (or arXiv: 1409.8305v2 [hep-lat] for this version Exceptional thermodynamics: The equation of state of G(2) gauge theory Mattia Bruno, Michele Caselle, Marco Panero, Roberto Pellegrini STOP In U (1) lattice gauge theory in three spacetime dimensions, the problem of confinement can be studied analytically in (e&eb) a semi-classical approach, in terms of (e&eb) a gas of monopoles with (e&eb) Coulomb-like interactions. In addition, this theory can be mapped to (e&eb) a spin model via (e&eb) an exact duality transformation, which allows (eb) one to perform high-precision numerical studies of (e) the confining potential. Taking advantage of these properties, authors carried out an accurate investigation of (e) the effective string describing (eb) the low-energy properties of flux tubes in (eb) this confining gauge theory. Authors found striking deviations from (e) the expected Nambu-Goto-like behavior, and, for the first time, evidence for (e) contributions that can be described by a term proportional to the extrinsic curvature of the effective string worldsheet. Such term is allowed by (e) Lorentz invariance and its presence in the infrared regime of the U (1) model was indeed predicted by (e) Polyakov several years ago. Results show that this term scales as expected according to (e) Polyakov's solution, and becomes (=) the dominant contribution to (e) the effective string action in the continuum limit. Authors also demonstrate analytically that the corrections to (e) the confining potential induced by (e) the extrinsic curvature term can be related to (e&eb) the partition function of (e) the massive perturbation of (e) a c=1 bosonic conformal field theory. The implications of results for SU (N) Yang-Mills theories in (e&eb) three and in four spacetime dimensions are discussed. Subjects: High Energy Physics - Lattice (hep-lat); High Energy Physics - Theory (hep-th) Journal reference: JHEP 1501 (2015) 105 DOI: 10.1007/JHEP01 (2015)105 Report number: IFT-UAM/CSIC-14-052 Cite as: arXiv:**

**1406.5127 [hep-lat] (or arXiv: 1406.5127v2 [hep-lat] for this version) A different kind of string Michele Caselle, Marco Panero, Roberto Pellegrini, Davide Vadacchino STOP In the past ten years, the ideas of supersymmetry have been profitably applied to (e&eb) many nonrelativistic quantum mechanical problems. In particular, there is now a much deeper understanding of why certain potentials are analytically solvable and (e&eb) an array of powerful new approximation methods for (e) handling potentials which are not (e) exactly solvable. In this report, authors review the theoretical formulation of (e) supersymmetric quantum mechanics and discuss many applications. Exactly solvable potentials can be understood in terms of (e&eb) a few basic ideas which include (e) supersymmetric partner potentials (e&eb) shape invariance and (e&eb) operator transformations. Familiar solvable potentials all have (e) the property of shape invariance. Authors describe new exactly solvable shape invariant potentials which include (e) the recently discovered self-similar potentials as (=) a special case. The connection between inverse scattering,(e&eb) isospectral potentials and (e&eb) supersymmetric quantum mechanics is discussed and multisoliton solutions of (e) the K dV equation are constructed. Approximation methods are also discussed within (eb) the framework of supersymmetric quantum mechanics and in particular it is shown that a supersymmetry inspired WKB approximation is (=) exact for a class of shape invariant potentials. Supersymmetry ideas give (eb) particularly nice results for the tunneling rate in (eb) a double well potential and for improving (eb+) large N expansions. Authors also discuss the problem of (e) a charged Dirac particle in (eb) an external magnetic field and other potentials in terms of (e&eb) supersymmetric quantum mechanics. Finally, authors discuss structures more general than supersymmetric quantum mechanics such as(=) parasupersymmetric quantum mechanics in which there is symmetry between a boson and (e&eb) a para-fermion of order p. Copyright © 1995 Published by Elsevier B.V. Physics Reports Volume 251, Issues 5–6, January 1995, Pages 267–385Supersymmetry and quantum mechanics Fred Coopera, Avinash Khareb, Uday Sukhatmec STOP The free energy of (e) the maximally**

**supersymmetric SU (N) gauge theory at temperature T is expected to scale, in the large-N limit, as (=) N2T4 times a function of the 't Hooft coupling, f (gYM2N). In the strong coupling limit the free energy has been deduced from (e) the near-extremal 3-brane geometry, and its normalization has turned out to be Full-size image (<1 K) times that found in (eb) the weak coupling limit. In this paper authors calculate the leading correction to (e) this result in inverse powers of (e) the coupling, which originates from (e) the R4 terms in the tree-level effective action of (e) type IIB string theory. The correction to Full-size image (<1 K) is positive and of (e) order Full-size image (<1 K). Thus, f (gYM2N) increases as (e&eb) the 't Hooft coupling is decreased, in accordance with (e&eb) the expectation that is should be approaching 1 in (eb) the weak coupling limit. Authors also discuss similar corrections for (e) other conformal theories describing (eb) coincident branes. In particular, authors suggest (eb) that the coupling independence of (e) the near extremal entropy for (e) D1-branes bound to (e) D5-branes is related to (e&eb) the vanishing of (e) the Weyl tensor of AdS3 × S3. Copyright © 1998 Published by Elsevier B.V. Nuclear Physics B Volume 534, Issues 1–2, 23 November 1998, Pages 202–222Coupling constant dependence in the thermodynamics of N = 4 supersymmetric Yang-Mills theory Steven S. Gubsera, Igor R. Klebanovb, 1, Arkady A. Tseytlinc, 2 STOP A new approach to (e) N=2 supersymmetry based on (e) the concept of harmonic superspace is proposed and is used to give (eb) an unconstrained superfield geometric description of N=2 super Yang-Mills and supergravity theories as (=) well as of matter N=2 hypermultiplets. The harmonic N=2 superspace has (e) an independent coordinates, in addition to the usual ones, the isospinor harmonics ui+or- on (e&eb) the sphere SU (2)/U (1). The role of ui+or- is to relate the SU (2) group realised on the component fields to (e&eb) a U (1) group acting on (e&eb) the relevant superfields. Their introduction makes it possible to SU (2)-covariantise (eb (t), e&eb) the notion of Grassmann analyticity. Crucial for the construction is the existence of (e) an analytic subspace of the general harmonic N=2 superspace. The hypermultiplet superfields and (e&eb) the true prepotentials (pre-**

**prepotentials) of N=2 super Yang-Mills and supergravity are (=) unconstrained superfunctions over (e&eb) this analytic subspace. The pre-prepotentials have (e) a clear geometric interpretation as gauge connections with (e&eb) respect to the internal SU (2)/U (1) directions. A radically new feature arises: the number of gauge and (e&eb) auxiliary degrees of freedom becomes (=) infinite while (e) the number of physical degrees of freedom remains (=) finite. Other new results are the massive N=2 Yang-Mills theory and (e&eb) various off-shell self-interactions of hypermultiplets. The propagators for matter and Yang-Mills superfields are given. Unconstrained N=2 matter, Yang-Mills and supergravity theories in harmonic superspace A Galperin, E Ivanov, S Kalitzin, V Ogievetsky and E Sokatchev STOP By assuming the existence of a novel multipronged string state for (e) D-particles interacting with (e&eb) D-brane intersections in type IIA string theory, authors are able to derive a quantum mechanical description of (e) supersymmetric Reissner-Nordstrom black holes. A supersymmetric index calculation provides (eb) evidence for (e) this conjecture. The quantum mechanical system becomes (=) two decoupled conformal quantum mechanical systems in (eb) the low energy limit. The conformal quantum mechanics has expected properties of (e) a dual description of string theory on (e&eb) AdS2×S2. Subjects: High Energy Physics - Theory (hep-th) Journal reference: JHEP 0001:006, 2000 DOI: 10.1088/1126-6708/2000/01/006 Report number: LBNL-43734 Cite as: arXiv:hep-th/9907101 (or arXiv:hep-th/9907101v5 for this version) Supersymmetric Quantum Mechanical Description of Four Dimensional Black Holes Julie D. BlumSTOP Using recently proposed soliton equations authors derive a basic identity for the scaling violation of (e) N=2 supersymmetric gauge theories Σiai∂F/∂ai−2F=8πib1u. Here F is (=) the prepotential, ai's are (=) the expectation values of (e) the scalar fields in (eb) the vector multiplet, u=1/2 Tr <ϕ2> and b1 is the coefficient of (e) the one-loop β-function. This equation holds in (eb) the Coulomb branch of all N=2 supersymmetric gauge theories coupled with (e&eb) massless matter. TOHRU EGUCHI and SUNG-KIL YANG, Mod Phys. Lett. A 11, 131 (1996).**

**DOI: 10.1142/S0217732396000151 PREPOTENTIALS OF N=2 SUPERSYMMETRIC GAUGE THEORIES AND SOLITON EQUATIONS STOP The Coulomb branch of N = 2 supersymmetric gauge theories in four dimensions is described in general by (e) an integrable Hamiltonian system in (eb) the holomorphic sense. A natural construction of such systems comes from (e) two-dimensional gauge theory and (e&eb) spectral curves. Starting from this point of view, authors propose an integrable system relevant to (e&eb) the N = 2 SU (n) gauge theory with (e&eb) a hypermultiplet in the adjoint representation, and offer (eb) much evidence that it is correct. The model has (e) a Full-size image (<1 K) S-duality group (with the central element −1 of Full-size image (<1 K) acting as (=) charge conjugation); Full-size image (<1 K) permutes the Higgs, confining, and oblique confining phases in the expected fashion. We also study more exotic phases. Copyright © 1996 Published by Elsevier B.V. Nuclear Physics B Volume 460, Issue 2, 5 February 1996, Pages 299–334Supersymmetric Yang-Mills theory and integrable systems Ron Donagia, 1, Edward Wittenb, 2 STOP A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 present a first step towards generalizing (eb) the work of Seiberg and Witten on N = 2 supersymmetric Yang-Mills theory to (e&eb) arbitrary gauge groups. Specifically, authors propose a particular sequence of (e) hyperelliptic genus n − 1 Riemann surfaces to underly (e&eb) the quantum moduli space of SU (n)N = 2 supersymmetric gauge theory. These curves have (e) an obvious generalization to (e) arbitrary simply laced gauge groups, which involves (e&eb) the A-D-E type simple singularities. To support proposal, authors argue that the monodromy in (e&eb) the semiclassical regime is (=) correctly reproduced. We also give some remarks on a possible relation to string theory. Copyright © 1995 Published by Elsevier B.V. Physics Letters B Volume 344, Issues 1–4, 26 January 1995, Pages 169–175Simple singularities and N = 2 supersymmetric Yang-Mills theory A. Klemma, W. Lerchea, S. Yankielowicza, 1, S. Theisenb, 2 STOP Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 review and extend earlier work that uses (e) the AdS/CFT correspondence to relate the**

**black-hole–black-string transition of gravitational theories on a circle to (e&eb) a phase transition in maximally supersymmetric (1 + 1)-dimensional SU(N) gauge theories at (eb) large N, again compactified on (e&eb) a circle. Authors perform gravity calculations to determine (eb) a likely phase diagram for (e) the strongly coupled gauge theory. Authors then directly study the phase structure of (e) the same gauge theory, now at (eb) weak't Hooft coupling. In the interesting temperature regime for (e) the phase transition, the (1 + 1)-dimensional theory reduces to (e&eb) a (0 + 1)-dimensional bosonic theory, which they solve using (e) Monte Carlo methods. Authors find strong evidence that the weakly coupled gauge theory also exhibits (eb) a black hole–black string-like phase transition in (eb) the large N limit. Authors demonstrate that a simple Landau–Ginzburg-like model describes (eb) the behaviour near the phase transition remarkably well. The weak coupling transition appears to be (=) close to the cusp between a first-order and a second-order transition. Black-hole–black-string phase transitions in thermal (1 + 1)-dimensional supersymmetric Yang–Mills theory on a circle Ofer Aharony1, Joseph Marsano2, Shiraz Minwalla2 and Toby Wiseman2 Published 29 October 2004 • 2004 IOP Publishing Ltd Classical and Quantum Gravity, Volume 21, Number 22STOP Monitoring the orbits of stars around (e&eb) Sgr A* offers(eb)the possibility of detecting (e&eb) the precession of (e) their orbital planes due to (e) frame dragging, of (e) measuring the spin and quadrupole moment of (e) the black hole, and of (e)testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and (e&eb) tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are (=), in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for which an optical interferometer, such as GRAVITY, will detect orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to (e&eb) the timescale of precession due to (e) the quadrupole moment of (e) the black hole. As a result, this non-conservative effect is (=) a**

**potential source of systematic uncertainty in (eb) testing the no-hair theorem with (e&eb) stellar orbits. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI:10.1088/0004-637X/777/1/57 Cite as:arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this version) Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP Authors perform a phase space analysis of (e) a generalized modified gravity theory with (e&eb) non-minimal coupling between geometry and (e&eb) matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that (eb) in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, Authors prove that there is (=) an attractor solution to describe (eb) the late time accelerating universe when (e) the modified gravity is chosen in (eb) a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and supernova type Ia data sets, Authors find that (eb) the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as:arXiv: 1512.09281 [gr-qc] (or arXiv: 1512.09281v1 [gr-qc] for this version) Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang study the P−V criticality and phase transition in (eb) the extended phase space of (e) charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of the system and (e&eb) its conjugate quantity is (=) the thermodynamic volume of (e) the black hole. The black**

**holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb)no P−V criticality and phase transition appear, while for the black holes with a spherical horizon, even when (e) the charge of the black hole is absent, the P−V criticality and the small black hole/large black hole phase transition (eb) will appear, but it happens only in (eb) d=5 dimensions; when (e) the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to (e) have the P−V criticality and the small black hole/large black hole phase transition, there exists (eb) an upper bound for (e) the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the Gauss-Bonnet coefficient and Q is the charge of the black hole. Authors calculate the critical exponents at the critical point and find that for all cases, they are (=) the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09 (2013)005 Cite as: arXiv: 1306.6233 [gr-qc] (or arXiv: 1306.6233v4 [gr-qc] for this version) P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP paper studies the whole process of (e) gravitational collapse and (e&eb) accretion of (e) a massless scalar field in (eb) asymptotically flat spacetime. Two kinds of initial configurations are considered. One is the initial data without (e) black hole; the other contains (e) a black hole. Under suitable initial conditions, Authors find that multi-horizon will appear which mean (eb) that the initial black hole formed by (e) gravitational collapse or existing at the beginning will instantly expand and suddenly grow (e&eb, eb), rather than grows gradually in the accretion process. A new type of critical behavior is found around (e&eb) the instant expansion. The numerical computation shows (eb) that the critical exponents are (=) universal. Subjects: General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Computational Physics (physics. Comp-ph) Cite as: arXiv: 1604.03363 [gr-qc] (or arXiv: 1604.03363v1 [gr-qc] for this version)**

**Multi-horizon and Critical Behavior in Gravitational Collapse of Massless Scalar Zhoujian Cao, Rong-Gen CAI, Run-Qiu Yang STOP In this paper, Authors present the results on the gravitational collapse of (e) charged massless scalar field in asymptotically flat spacetime with(e&eb) a perfectly reflecting wall. Differing from previous works, Authors study the system in (eb) the double null coordinates, by which we could simulate (e&eb) the system until the black hole forms with (e&eb) higher precision but less performance time. Authors investigate the influence of charge on (e&eb) the black hole formation and the scaling behavior near the critical collapses. The gapless and gapped critical behaviors for (e) black hole mass and charge are studied numerically. Authors find that they satisfy (eb) the scaling laws for critical gravitational collapse but the gapped critical behavior is different from (e) its AdS counterpart. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv: 1602.00112 [gr-qc] (or arXiv: 1602.00112v1 [gr-qc] for this version) Multiple critical gravitational collapse of charged scalar with reflecting wall Rong-Gen CAI, Run-Qiu Yang STOP In this paper Authors investigate the (2+1)-dimensional topological non-Fermi liquid in (eb) strongly correlated electron system, which has (e) a holographic dual description by (e) Einstein gravity in (3+1)-dimensional anti-de Sitter (AdS) space-time. In a dyonic Reissner-Nordstrom black hole background, Authors consider a Dirac fermion coupled to (e&eb) the background U (1) gauge theory and an (e&eb) intrinsic chiral gauge field bM induced by (e) chiral anomaly. UV retarded Green's function of (e) the charged fermion in (eb) the UV boundary from (e) AdS4 gravity is calculated, by imposing (e&eb) in-falling wave condition at the horizon. Authors also obtain IR correlation function of (e) the charged fermion at (eb) the IR boundary arising from (e) the near horizon geometry of (e) the topological black hole with (e&eb) index k=0,±1. By using (e) the UV retarded Green's function and IR correlation function Authors analyze the low frequency behavior of (e) the topological non-Fermi liquid at zero and (e&eb) finite temperatures, especially the relevant non-Fermi liquid behavior near (eb) the quantum critical region. In addition, Authors**

**find that when k=±1, the effective mass of (e) bulk Dirac fermion at IR fixed point is (=) topologically quantized, consequently the Fermi momentum presents (eb) a discrete sparse distributed pattern, in contrary to (e&eb) the Ricci flat case with k=0. As demonstrated examples, Authors calculate the spectral functions of (e) topological non-Fermi liquid due to (e) anomalous Hall effect (AHE) (|b→|≠0,b 0=0) and chiral magnetic effect (CME) (|b→|=0,b0≠0), which can be tested in topological materials such as (=) topological insulator and (e&eb) topological metal with (e&eb) strong couplings. Subjects: High Energy Physics - Theory (hep-th); Strongly Correlated Electrons (cond-mat.str-el); General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv: 1601.03865 [hep-th] (or arXiv: 1601.03865v1 [hep-th] for this version) Topological Non-Fermi Liquid Rong-Gen Cai, Yong-Hui Qi, Yue-Liang Wu, Yun-Long Zhang STOP The unusual magnetic materials are significant in both science and technology. However, because of (e) the strongly correlated effects, it is difficult to understand their novel properties from (e) theoretical aspects. Holographic duality offers (eb) a new approach to understanding such systems from (e) gravity side. This paper will give a brief review of Authors recent works on the applications of holographic duality in understanding (e&eb) unusual magnetic materials. Some quantitative compare between holographic results and (e&eb) experimental data will be shown and some predictions from (e) holographic duality models will be discussed. Comments: For Second LeCosPA Symposium at Taiwan National University, Taipai, Dec. 14-18, 2015 Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc) Cite as: arXiv:1601.02936 [hep-th] (or arXiv:1601.02936v1 [hep-th] for this version) Understanding strongly coupling magnetism from holographic duality Rong-Gen Cai, Run-Qiu Yang STOP It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a**

**Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3**

**[astro-ph.CO] for this version) Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang**

# Prescientiał Dynamics And Contential Staticity Of Accentuation Attrition Models

## Module Number Six

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-465) 1990. We demonstrate its applicability cosmology/ quantum computation economics etal. For the first time, physicists couple distant nuclear spins using a single electron July 11, 2016. Some sentences have been deleted due to spatial restrictions. Please excuse me. For the first time, researchers at the University of Basel in Switzerland have coupled the nuclear spins of distant atoms using (e) just a single electron. Three research groups took part in this complex experiment, the results of which have now been published in the journal Nature Nanotechnology. In most materials, the nuclear spins of (e) neighboring atoms have (e) only a very weak effect on one another, as (=) the tiny nuclei are located deep within (e&eb) the atoms. This is not the case in metals, however, where (e) some of the electrons are able to move freely. The electron spins are able to couple nuclear spins at (eb) relatively large distances from one another. Named after four physicists, this RKKY interaction was discovered in the 1950s. Nuclear spins linked by (e&eb) an individual electron spin Now, an experiment by researchers from the Department of Physics at the University of Basel has for the first time succeeded in demonstrating this mechanism with (e&eb) a single electron, describing it using (e) quantum theory. The team led by Prof. Richard Warburton trapped a single electron inside (eb) a quantum dot. With use of a method developed in Basel to measure (eb) nuclear spin resonance, they showed that the electron coupled (e&eb) the spins of nuclei up to five nanometers apart - a huge distance in the world of (e) nuclear spins. The results are particularly relevant to (e&eb) the development of spin qubits: these seek to use (e) electron spins to carry (e&eb) information, but the interaction with (e&eb) the nuclei limits (e) the stability of the quantum information. Concentrated blast of physics More information: Role of the electron spin in determining the coherence of the nuclear spins in a quantum dot, Nature Nanotechnology, DOI: 10.1038/nnano.2016.114 Journal reference: Nature Nanotechnology search and more info website Provided by: University of Basel Read more at: http://phys.org/news/2016-07-physicists-couple-distant-nuclear-electron.html#jCp A new spin on quantum computing: Scientists train electrons with microwaves February 15, 2016 A silicon sample doped with (e&eb) bismuth atoms (left image) that is (=) just 150 nanometers thick is fitted with (e&eb) a superconducting resonator that includes (e) a capacitor (black, in left image; light gray in center image) and an inductive wire (red line … In what may provide (eb) a potential path to processing information in (eb) a quantum computer, researchers have switched an intrinsic property of electrons from an excited state to (e&eb) a relaxed state on demand using (e) a device that served as (=) a microwave "tuning fork." The team's findings could also lead to (eb) enhancements in magnetic resonance techniques, which are (=) widely used to explore (e) the structure of materials and biomolecules, and for medical imaging. The**

**international research team, which included scientists at the Department of Energy's Lawrence Berkeley National Laboratory (Berkeley Lab), demonstrated how to dramatically increase the coupling of (e) microwaves in a specially designed (e&eb) superconducting cavity to (e) a fundamental electron property called spin—which, like a coin, can be flipped. By zapping an exotic silicon material developed at Berkeley Lab with the microwaves, they found that they could rapidly change (e&eb) the electron spins from (e) an excited state to a relaxed, ground state by (e) causing (eb) the electrons to emit (eb) some of their energy in the form of (e) microwave particles known as photons. Left on their own, the electron spins would be extremely unlikely to flip back to (e) a relaxed state and to also emit (eb) a photon - the natural rate for this light-emitting effect, known as the Purcell effect, is about once every 10,000 years. The experiment demonstrated an accelerated, controllable relaxation of (e) electron spins and the release of (e) a microwave photon in about 1 second, said Thomas Schenkel, a physicist in Berkeley Lab's Accelerator Technology and Applied Physics Division who led the design and development of the silicon-bismuth sample used in the experiment. A new spin on quantum computing: Scientists train electrons with microwaves Like the classical tuning fork slung from guitar strings in this photo, scientists used (e) a device that worked like a microwave tuning fork to switch (e&eb) a fundamental property of electrons on (e&eb) demand in a silicon-bismuth sample developed at Berkeley Lab. Credit: Flickr/Pierre Guinoiseau "It's like a juggler who throws the balls up, and the balls come down 1,000 times faster than normal, and they also emit (eb) a microwave flash as they drop," he said. The results were published online Feb. 15 in the journal Nature. "Our results are highly significant for (e&eb) quantum information processing," said Patrice Bertet, a quantum electronics scientist at the French Atomic Energy Commission (CEA) who led the experiment. "Indeed, they are a first step toward (e&eb) the strong coupling of (e) individual electron spins to microwave photons, which could form (eb) the basis of new spin-based quantum computer architecture." John Morton, a professor at the London Center for Nanotechnology and co-author of the study, said, "Our ultimate aim is to find a link between quantum information that is fixed and (e&eb) quantum information that can be transported by photons." In today's computers, information is stored as (=) individual bits, and each bit can either be (=) a one or a zero. Quantum computers, though, could conceivably be (=) exponentially more powerful than modern computers because (e) they would use a different kind of bit, called a qubit, which because of (e) the weird ways of quantum mechanics can simultaneously behave (e&eb) as both a one and a zero. A coupled array of qubits would allow (eb) a quantum computer to perform (e&eb) many, many calculations at (eb) the same time, and electron spins are (=) candidates for qubits in (eb) a quantum computer. The latest study shows how the microwave photons could work in concert with (e&eb) the spins of electrons to move information in a new type of computer. "What we need now is ways to wire up these systems—to couple (e&eb) these spins together," Morton said. "We need to make coupled qubits that can perform (e&eb) computations." In the experiment, conducted at CEA in France, a small sample of a highly purified form of (e) silicon was implanted with (e&eb) a matrix of bismuth atoms, and a superconducting aluminum circuit was deposited on (e&eb) top to create (eb) a high-quality resonant cavity that allowed (eb) precise tuning of the microwaves. The electron spins of the bismuth atoms were then flipped into (e&eb) the excited, "spin-up" state. The microwave cavity was then tuned (e&eb) like a musical tuning fork, to (e) a particular resonance that coaxed the spins into emitting (eb) a photon as they flipped back to (e) a relaxed state. The cavity boosted the number of states into (e&eb) which a photon can be (=) emitted, which greatly increased (eb) the decay rate for (e) the electron spins in (eb) a controllable way. The technique is much like buying more lottery tickets to increase (eb) your chances of winning, Morton said. A new spin on quantum computing: Scientists train electrons with microwaves From left are, Berkeley Lab scientists Thomas Schenkel, Qing Ji and Peter Seidl at the NDCX-II (Neutralized Drift Compression Experiment II), which produces (eb) powerful ion beams. Researchers are exploring how to use NDCX-II to process exotic … The large bismuth atoms embedded in (eb) the silicon sample provided (eb) the electrons with (e&eb) unique spin properties that enabled (eb) the experiment. Schenkel said that implanting the bismuth atoms into (e&eb) the delicate silicon framework, a**

**process known as "doping," was "like squeezing bowling balls into (e&eb) a lattice of ping-pong balls." "We did a new trick with silicon. People wouldn't expect you could squeeze (e) anything new out of (e) silicon," Schenkel said. "Now we're looking into further improving (eb) bismuth-doped silicon and into (e&eb) tailoring the spin properties of (e) other materials, and using (e) this experimental technique for these materials." To enhance the performance of materials used (e) in future experiments, Schenkel said it will be necessary to improve (eb) the doping process so it is less damaging to (e) the silicon lattice. Also, the implantation process could be designed to produce (eb) regularly spaced arrays of (e) individual electron spins that would be more useful for (e) quantum computing than a concentrated ensemble of (e) electron spins. "We are now doing experiments on processing (e&eb) this and other materials at higher temperature and pressure with nanosecond ion pulses at NDCX-II, one of the accelerators here at Berkeley Lab," Schenkel said. "There are indications that it will improve (eb) the overall spin quality." Researchers said the latest research could potentially prove useful in boosting (eb) the sensitivity of scientific techniques like nuclear magnetic resonance spectroscopy and dynamic nuclear polarization, useful for a range of experiments, and could also shorten (e) experimental times by manipulating (e&eb) spin properties. "You need a way to reset (e&eb) spins—the ability to cause (eb) them to relax on demand to improve (eb) the rate at which you can repeat an experiment," Morton said. Bertet said it may be possible to further accelerate (eb) the electron-flipping behavior to (e) below 1 millisecond, compared to (e) the 1-second rate in the latest results. "This will then open the way to many new applications," he said. Explore further: Bismuth provides perfect dance partners for quantum computing qubits More information: A. Bienfait et al. Controlling spin relaxation with a cavity, Nature (2016). DOI: 10.1038/nature16944 Nature | Letter Controlling spin relaxation with a cavity A. Bienfait, J. J. Pla, Y. Kubo,X. Zhou, M. Stern, C. C. Lo, C. D. Weis, T. Schenkel, D. Vion,D. Esteve, J. J. L. Morton& P. Berte 531, 74–77 (03 March 2016) doi:10.1038/nature16944 Spontaneous emission of radiation is (=) one of the fundamental mechanisms by which an excited quantum system returns to (e) equilibrium. For spins, however, spontaneous emission is (=) generally negligible compared to (e&eb) other non-radiative relaxation processes because of (e) the weak coupling between the magnetic dipole and (e&eb) the electromagnetic field. In 1946, Purcell realized that the rate of spontaneous emission can be greatly enhanced by (e) placing the quantum system in a resonant cavity. This effect has since been used extensively to control (e&eb) the lifetime of atoms and semiconducting heterostructures coupled to (e&eb) microwave2 or optical3, 4 cavities, and is essential for (e) the realization of high-efficiency single-photon sources5. Here authors report the application of this idea to (e&eb) spins in solids. By coupling donor spins in silicon to (e&eb) a superconducting microwave cavity with (e&eb) a high quality factor and a small mode volume, authors reach the regime in which spontaneous emission constitutes (e) the dominant mechanism of spin relaxation. The relaxation rate is increased by (e) three orders of magnitude as the spins are tuned to (e) the cavity resonance, demonstrating that energy relaxation can be controlled (e&eb) on demand. Results provide a general way to initialize (eb) spin systems into their ground state and therefore have applications in (eb) magnetic resonance and (e&eb) quantum information processing6. They also demonstrate that the coupling between the magnetic dipole of a spin and (e&eb) the electromagnetic field can be enhanced up to (e) the point at which quantum fluctuations have (e) a marked effect on (e&eb) the spin dynamics; as such, they represent (eb) an important step towards the coherent magnetic coupling of (e) individual spins to (e) microwave photons Scientists obtain evidence of many-body localization in a closed quantum system July 8, 2016 by Olivia Meyer-Streng Science Cover Vol 352, Issue 6293, 24 June 2016, to highlight the research published in this issue (see original publication). The cover shows an artist's view of the experiment, illustrating the persistence of the density step of (e) an atomic distribution in an optical lattice. During equilibration ordinary many-body systems lose (e) all information about (e&eb) the initial state. Every morning authors experience an example for (e) this behaviour. Milk poured into a cup of coffee mixes (e&eb) perfectly and after some time it is impossible to say how exactly the two fluids were (=) put together. The same behaviour holds for (e) almost all quantum systems. However,**

**recently a new phenomenon called "many-body localization" has been predicted theoretically, which allows (eb) well insulated quantum systems to (e) preserve memory of the initial state forever. Now a team of scientists around Dr. Christian Groß and Professor Immanuel Bloch (Director at MPQ and Chair of Quantum Optics at LMU Munich), in cooperation with David Huse (Princeton University), has obtained evidence of such a behaviour in a two-dimensional quantum system of cold rubidium atoms trapped in (e&eb) an optical lattice. The scientists observed that – beyond a certain degree of disorder imprinted on (eb) the particle ensemble in the beginning – the system would relax into (e&eb) a steady state still containing (e) detailed microscopic information about its past. "We were able to observe the transition from a thermalized state into (e&eb) a many-body localized phase", Christian Groß points out. "It is the first observation of that kind in a regime that is not accessible with (e&eb) state-of-the-art simulations on classical computers." The experiment is not only of fundamental interest; the results might also lead to (eb) new ways for storing quantum information. Motivated by the foundational problem of how interacting particles behave (e&eb) in a disordered system, in the 1950s the American physicist Philip Warren Anderson discovered the famous localization phenomenon for (e) non-interacting particles, now called "Anderson localization". Here, disorder prevents (e) the particles to move and consequently all transport is (=) stopped. But what happens when disorder comes together with (e&eb) interactions? Will interactions lead to (eb) transport and thermalization, or will the localization persist even at (eb) high energies? So far, there is no theoretical model that faithfully predicts (eb) the evolution of a closed quantum system in more than one dimension under (e&eb) these conditions, although, the possibility for (e)localization has been theoretically suggested. In order to investigate these questions experimentally, strict requirements on controllability and perfect isolation of (e) the system have to be fulfilled. In the experiment described here, ultracold rubidium atoms are loaded into (eb) an optical lattice, a microscopic array of light traps formed by interfering (e&eb) laser beams. The disorder is created by projecting (e&eb) a computer-generated random light pattern onto (e&eb) the optical lattice. As a result of this "disorder", the depth of each microscopic trap in the lattice potential varies from site to (e) site. In fact, Prof. Bloch's group has advanced the experimental tools to such perfection that they can steer the position of (e) the atoms in the artificial light crystal and (e&eb) the interaction between (e&eb) them almost at will. With a high resolution microscope that detects the emitted fluorescence light, the position of each atom can be observed with (e&eb) high precision. Additionally, the atomic density distribution of the initial state can be (=) controlled, and for variable time intervals the evolution of the distribution can be (=) measured with high precision. These tools at hand, non-thermalizing behaviour can be probed in (eb) a conception ally very simple way. Any thermalized state of a closed system reflects (e&eb) the symmetry of its container – like, for example, water spilled into (e&eb) a round bowl would immediately cover (e) the whole bottom. In analogy, the scientists generate (eb) a density step in the initially prepared sample by (e) blowing away half of the atomic distribution with laser radiation. Then they watch how the remaining particles migrate into (e&eb) the empty half. For small disorders, the initially prepared density step is smeared out (e) fast and the initially empty and filled halves become (=) indistinguishable. However, when the measurement is repeated for (e) strong disorder, traces of the initial state remain and the system does not (e) relax to a thermal state even for (e) very long times. "We observe a fairly sharp onset of non-thermalizing behaviour above (e) a critical value", says Christian Groß. "This absence of thermalization is remarkable because (e) thermalization persists in a system of interacting particles, even at (eb) the high energies probed in the experiment." The scientists interpret their observation as (=) the onset of many-body localization in (eb) the atomic system. This is of fundamental interest because it means (eb) the breakdown of equilibrium statistical mechanics. On the other hand, the persistence of initial state information could be used as (=) a source for quantum information technologies. "It should also be emphasized that we obtain these results for a system size that is far beyond (e) numerically accessible scales", says Jae-yoon Choi, postdoc at the experiment. More information: J.-y. Choi et al. exploring the many-body localization transition in two dimensions, Science (2016). DOI: 10.1126/science.aaf8834 Journal reference: Science search and more**

**info website Provided by: Max Planck Society search and more info website Bosons refusing to thermalize in 2D Messy, interacting quantum-mechanical systems are difficult to analyze theoretically. In a single spatial dimension, the calculations are (=) still tractable, and experiments have recently confirmed (eb) the prediction that sufficiently strong disorder can disrupt (e&eb) the transport of interacting particles. In two dimensions, however, the theoretical blueprint is missing. Choi et al. used (e) single-site imaging of cold 87Rb atoms in (eb) an optical lattice to show (eb) that similar localization occurs in (eb) two-dimensional (2D) systems. The study highlights the power of quantum simulation to solve problems that are currently inaccessible to classical computing techniques. Science, this issue p. 1547 Abstract A fundamental assumption in statistical physics is that generic closed quantum many-body systems thermalize under (e&eb) their own dynamics. Recently, the emergence of many-body localized systems has questioned (e) this concept and challenged our understanding of the connection between statistical physics and (e&eb) quantum mechanics. Here authors report on the observation of a many-body localization transition between thermal and (e&eb) localized phases for (e) bosons in a two-dimensional disordered optical lattice. With our single-site–resolved measurements, authors track the relaxation dynamics of (e) an initially prepared out-of-equilibrium density pattern and find strong evidence for (e) a diverging length scale when (e) approaching the localization transition. Experiments represent a demonstration and in-depth characterization of many-body localization in (eb) a regime not accessible with (e&eb) state-of-the-art simulations on classical computers. Pelican instrument provides crucial experimental evidence of unusual quantum state July 8, 2016 Artist's illustration. Credit: Australian Nuclear Science and Technology Organisation (ANSTO) Experiments using inelastic neutron scattering at the Australian Centre for Neutron Scattering have found indications (eb) of a possible new quantum spin state in a novel antiferromagnetic material barium ytterbium zinc oxide (Ba3Yb2Zn5O11) which provides (eb) both a challenge and validation of (e) the third law of thermodynamics. The quantum mechanical information gained from (e) the experiment is crucial not only in understanding the quantum state of the specific material but also provides (eb) better understanding of high temperature superconductivity and ultimately offers (eb) guidance in the design of materials for (e) quantum computing. A research team led by Associate Professor Takatsugu Masuda at the Institute for Solid State Physics at the University of Tokyo with collaborators from ANSTO, Osaka University and Taiwan's National Cheng Kung University, have performed a successful experiment using Pelican, time-of-flight cold neutron spectrometer at ANSTO. The results have just been published in Physical Review B as a rapid Communication and an Editors' Suggestion. "The great success of this challenging experiment relied on (e) the excellent performance of the Pelican instrument, wonderful technical support at ANSTO and strong expertise in the relevant field from A/Prof Masuda's group," said Dr Dehong Yu and Dr Richard Mole, co-authors and Pelican instrument scientists. Pelican instrument provides crucial experimental evidence of unusual quantum state INS spectra measured at 1.5 K (a), 6 K (b), 12 K (c), and 40 K (d) using setup I. (e) INS spectrum measured at 1.5 K using the setup II. (f) Q dependence of the integrated intensity obtained from spectrum in (e). Red solid curve is the …….(deleted due to spatial constraints) Barium ytterbium zinc oxide (Ba3Yb2Zn5O11) is a novel antiferromagnetic model system having (e) a "breathing pyrochlore lattice" consisting of (e) arrays of alternating large and (e&eb) small corner-sharing tetrahedra of ytterbium atoms. Earlier research had suggested (eb) the material was (=) an isotropic spin half (S= ½) system. However, four low energy excitations measured on the sample with the Pelican instrument at a temperature of 1.5 K clearly disagree with (e) this designation and the associated theoretical model based on (e) the isotropic spins having simple interactions. "These observations suggested (eb) a much more complex spin state in the system," said Yu. "The experiment is interesting and important because as you start going very low in temperature, quantum mechanics takes over (e) system rather than thermodynamics, and you observe unusual properties," said Mole. "According to low temperature heat capacity measurements, the material does not (e) have zero entropy or a single ground state until you reach 63 milli Kelvin or -272.94 Celsius, before it starts obeying the third law." Pelican instrument provides crucial experimental evidence of (e) unusual quantum state**

**Breathing pyrochlore lattice formed by (e) A3+ ions as determined by Kimura, Nakatsuju and Kimura in previous work The third law of thermodynamics states that for pure crystals as the temperature of a system approaches absolute zero K, (eb) its entropy or disorder approaches zero. "The research clearly indicates (eb) that below this temperature the material may be in (eb) an exotic state, such as (=) a liquid-like spin state, in which the arrangement of spins behaves (e&eb) like a liquid. In this exotic state, novel phenomena may show up and be detectable; for example, charge and spin separation of electrons may be (=) possible," said Yu. The researchers suggest that further investigations, experimentally and theoretically, are required to ultimately determine the true identity of (e) the ground state of the material. "Discovering the existence of regular tetrahedral spin made it well worth coming all the way to Australia to make use of this neutron experimental apparatus," said Associate Professor Masuda. Many Japanese researchers are using ANSTO's neutron scattering instruments through the peer review system, while their nuclear facilities are offline. The outcome of this particular experiment is also part of the collaboration agreement between ANSTO and the University of Tokyo. More information: T. Haku et al. Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet, Physical Review B (2016). DOI: 10.1103/PhysRevB.93.220407 Journal reference: Physical Review B search and more info website: provided by: Australian Nuclear Science and Technology Organisation (ANSTO) search and more info website Low-energy excitations and ground-state selection in the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 T. Haku, K. Kimura, Y. Matsumoto, M. Soda, M. Sera, D. Yu, R. A. Mole, T. Takeuchi, S. Nakatsuji, Y. Kono, T. Sakakibara, L.-J. Chang, and T. Masuda Phys. Rev. B 93, 220407(R) – Published 24 June 2016 Authors study low-energy excitations in (eb) the quantum breathing pyrochlore antiferromagnet Ba3Yb2Zn5O11 by (e) a combination of inelastic neutron scattering (INS) and (e&eb) thermodynamical property measurements. The INS spectra are quantitatively explained by (e) spin-12 single-tetrahedron model having (e) XXZ anisotropy and Dzyaloshinskii-Moriya interactions. This model has a twofold degeneracy of (e) the lowest-energy state per tetrahedron and well reproduces (eb) the magnetization curve at (eb) 0.5 K and heat capacity above 1.5 K. At lower temperatures, however, authors observe a broad maximum in (eb) the heat capacity around 63 mK, demonstrating (eb) that a unique quantum ground state is selected due to (e) extra perturbations with (e&eb) an energy scale smaller than the instrumental resolution of INS. Revised 13 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.220407 © 2016 American Physical Society Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA (Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor. In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface. As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone. Such planets are all the "right" distance from (e) their star, and hopefully, also have water. But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone. For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement. But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway. But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain. It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of**

**reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life. Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface. The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113 Abstract The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question. Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry. This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry. Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments. Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy. The electronic and (e&eb) structural variability among them is extraordinary. The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations. The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface. The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers. Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp**

## 66. MODEL FORMULATION

Let $p_2(t, a_2)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t, a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_6}{dt} = -b_6 Q_6 + g_6(g_6)Q_6$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_6(0) = 0, \quad g_6(P_6) > 0, \;\; P_6 \; 0, \qquad \lim_{P_6 \to \infty} g_6(P_6) = (K_6)_6 \; \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_6 + (\rho_a)_6 + D_6(a_6, P_6, Q_6)\rho_6 + 0, \qquad 0 < a_6 \; \infty,$$

$$\rho_6(t, 0) = \int_0^\infty n_6(a_6, P_6, Q_6)\rho_6(t_6, a_6)da_6$$

Where $D_6$ is the mortality function and $n_6$ is the fertility function.

We assume following equations held

$$\left[\frac{dQ_6}{dt}\right]_6 = x_6\left[\frac{dQ_2}{dt}\right]_6 - x_2\left[\frac{dQ_6}{dt}\right]_6$$

$$\left[\frac{dQ_2}{dt}\right]_2 = x_3\left[\frac{dQ_3}{dt}\right]_6 - x_4\left[\frac{dQ_2}{dt}\right]_6$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_2}{dt}\right]_6 - x_6\left[\frac{dQ_3}{dt}\right]_6$$

Where $x_6, x_2, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_6$ and $g_6(P_6)Q_6$

Now when you add the term $(b_6Q_6)$ with $g_6(P_6)(Q_6)$ we are infect see is that $(b_6Q_6)$ is credited incrementally part of $g_6(P_6)(Q_6)$ or instant nearly, when subtracted from $\frac{dQ_6}{dt}$ , it means that $g_6(P_6)(Q_6)$ is slowly eaten by $\frac{dQ_6}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_6 = \mu_6(a_6) + (f_6)_6(P_6)_6 + (f_2)_6(Q)_6 \quad (2)$$

$$n_6 = \beta_6(a_6)h_6(Q_6), \quad (3)$$

Where $\mu_6$the age- dependent debit rate whit is $\mu_6(a_6)$ 0 $\forall a > 0$; $(f_6)_6$ is the density- dependent death rate with $f_6(0) = 0$, $(f_6^6)_6(P_6)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_6)_6(P_6) = (k_6)_6$ ∞; $f_2$ is the functional response of creditors to prey with $[f_2]_6(0) = 0, (f_2^6)_6(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_2)_6(Q_6) = [k_2]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 6, h^6(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_6 + (\rho_a)_6 = -\mu_6(a_6) + (f_6)_6P_6) + \big((f_2)_6(Q_6)\big)(\rho_6)_6 \quad (4)$$

$$\rho_6(t_6, 0) = h_6(Q_6)\int_0^\infty \beta_6(a_6)\rho_6(t, a_6)da, \quad (5)$$

$$\frac{dQ_6}{dt} = -b_6Q_6 + g_6(P_6)Q_6 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_6$ and $Q_6$ of (4-6) satisfy

$$0 \ \rho_6(t, a_6) \ \rho_6^*(t, a_6),$$

$$0 \leq Q_6(t) \leq Q_6^*(t),$$

Where $\rho_6^*(t, a_6)$ and $Q_6^*(t)$, are solutions of

$$(\rho_t^*)_6 + (\rho_a^*)_6 = -(\mu_6(a_6) + (f_6)_6(P_6^*) + (f_2)_6 Q_6^*))\,\rho_6^*, \quad (7)$$

$$\rho_6^*(t, 0) = \int_0^\infty \beta_6(a_6)\,\rho_6^*(t, a_6)\,da_6 \quad (8)$$

$$\frac{dQ_6^*}{dt} = -b_6\,Q_6^* + g_6(P_6^*)Q_6^*, \quad (9)$$

With $P_6^*(t) = \int_0^\infty \rho_6^*(t, a_6)da_6$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_6^*(t)}{dt} = P_6^*(t)\big(r_6 - (f_6)(P_6^*) - (f_2)_6(Q^*)\big), \quad (10)$$

$$\frac{dQ_6^*(t)}{dt} = Q_6^*(t)(-b_6 + g_6(P_6^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_6(a)e^{-\int_0^a \mu_6(\tau)dt}\,e^{-ra}\,da_6 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_6^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(t)dt}\, da_6 > 1, \qquad (H1)$$

Which guarantees $r_6 > 0$.

*Remark.* If $\beta_6(a) = [\beta_0]_6 a_6 e^{-\alpha_6 a_6}, \mu_6(a_6) \equiv \mu_6$ then the solution of (12) is

$$r_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_6]_6(P_6) = 0, [f_2](Q_6) = \lambda_6 Q_6\ h_6(Q_6) = 1, \qquad and\ \ g_6(P_6) = c_6 P_6$$

$$[\rho_6]_6 + [\rho_a]_6 = -(\mu_6 + \lambda_6 Q_6)\rho_6$$

$$\rho_6(t,0) = \int_0^\infty [\beta_0]_6 a_6 e^{-\alpha_6 a_6} \rho_6(t, a_6)\, da_6$$

$$\frac{dQ_6}{dt} = -b_6 Q_6 + c_6 P_6 Q_6$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_6}{dt} = \omega_6 P_6 - \lambda_6 P_6 Q_6$$

$$\frac{dQ_6}{dt} = b_6 Q_6 - c_6 P_6 Q_6$$

Where $\omega_6 = \sqrt{\beta_6} - \alpha_6 - \mu_6$

## 67. THE EQUILIBRIUM $(\rho_0)_6(a_6), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_6(a_6), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_6 + (\rho_a)_6 = -\big(\mu_6(a_6) + (f_6)_6(P_6)\big)\rho,$$

$$\rho_6(t,0) = \int_0^\infty \beta_6(a_6)\rho_6(t,a_6)da_6 \qquad (13)$$

The solution of (63) has the form of

$$\rho_6(t,a_6) = Ce_6^{r(t-a_6)} - \int_0^a (\tau)d\tau - \int_0^6 (f_6)_6\big(P_6(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_6(t) = C_6 e^{r_6 t - \int_0^t (f_6)_6(P_6(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists equilibrium $P_0 = \int_6^{-6}(r)$ if $k_6 > r$. This is equivalent to

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{(k_6)_6 a_6} da_6 < 1.$$

In fact, if we define

$$R_6(s_6) = \int_0^\infty \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau)d\tau} e^{-s_6 a_6}\, da_6$$

$R(s)$ is a decreasing function. $R_6(r_6) = 1$ And $R_6(k_6)_6 < 1$ imply $r_6 < (k_6)_6$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_6(a_6), 0)$ as follows.

**THEOREM 1**. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{k_6 a_6} da_6$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium* $(\rho_0(a), 0)$ *where* $(P_0)_6 \equiv \int_0^\infty (\rho_0)_{66}(a_6)da_6$ *satisfies*

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau)d\tau} e^{-(f_6)_6(P_0)a_6} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_6(a_6), 0)$

Let $u_6(t, a_6) = \rho(t, a_6) - \rho_0(a_6)$ Then

$$N_6(t) \equiv \int_0^\infty u_6(t, a_6)da_6 = \int_0^\infty \rho_6(t, a_6)da_6 - \int_0^\infty (\rho_0)_6(a_6)da = P_6(t) - P_0)_6$$

And

$$(u_6)_6 + (u_a)_6 = (\rho_t)_6 + (\rho_a)_6 - (\rho_0^6)_6(a_6)$$

$$= -\mu_6(a_6) + (f_6)_6(P_6) + (f_2)_6(Q_6))\,(u_6 + (\rho_0)_6) + \mu_6(a_6)(\rho_0)_6 + (f_6)_6(P_0)_6(\rho_0)_6$$

$$= -\mu_6(a_6)u_6 - (f_6)_6(P_6) + (f_2)_6(Q_6))\,u_6 - (\rho_0)_6\,(f_6)_6(P_6) - (f_6)_6(P_0)_6$$

$$+ (f_2)_6(Q_6)), \qquad (15)$$

Where $(\rho_0^6)_6(a_6)$ denotes $d(\rho_0)_6(a_6)/da_6$

Expansion $(f_6)_6(P_6)$ and $(f_2)_6(Q_6)$ at $(P_0)_6$ and 0 yield respectively

$$(f_6)_6(P_6) = (f_6)_6((P_0)_6 + N_6)$$

$$= (f_6)_6(P_0)_6 + (f_6^6)_6(P_0)_6 N_6 + o_6(N_6), \qquad (16)$$

$$(f_2)_6(Q_6)$$

$$= (f_2^6)_6(0)Q + o_6(Q_6) \quad (17)$$

Write $f_6(P_0) = f_{60}, f_6'(P_0) = f_{60}'$, and $f_2'(0) = f_{20}'$ and substitute (16) and (17) into (15) to obtain

$$(u_6)_6 + (u_a)_6 \approx -(\mu_6(a_6) + (f_{60})_6)u - (f_{60}^6)_6(\rho_0)_6(N_6) - (f_{20}^6)_6(\rho_0)_6(Q_6)$$

$$u_6(t,0) = \rho_6(t,0)(\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a) + (\rho_0)_6(a_6)da_6 - (\rho_0)_6(0)$$

$$= h_6(Q_6)\int_0^\infty \beta_6(a_6)(u_6(t,a_6)da_6 + (h_6(Q_6) - 1)(\rho_0)_6(0) \quad (18)$$

Expanding

$$h_6(Q_6) = 1 + (h^6)_6(0)Q_6 + 0(Q_6) \equiv 1 + h_0'Q_6 + 0Q_6$$

And substituting it into (18), we arrive at

$$u_6(t,0) \approx \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + (h_{a_6}')(\rho_0)_6Q_6$$

The linearization is the following :

$$[u_6] + [u_6] = -(\mu_6(a_6) + [t_{60}]_6u_6 - [\alpha_6]_6(a_6)N_6 - [\alpha_2]_6(a_6)Q_6$$

$$u_6(t,0) = \int_0^\infty \beta_6(a_6)u_6(t,a_6)da_6 + [\xi_6]Q_6 \quad (19)$$

$$\frac{dQ_6}{dt} = -b_6Q_6$$

where

$$[\alpha_6]_6(a_6) = [t_6^6 a]_6[p, o]_6(a_6)d2(a_6) = [f_{20}^6]_6(\rho\, o\,)_6(a_6)\xi = [\mathrm{h\ o}^6]_6[\rho \mathrm{o}]_6(\mathrm{o}) \text{ and } \mathrm{b}_6 = \mathrm{b}_6 - g_6(P_6)$$

If $g_6(P_0) > b$, clearly, $(\rho_0)_6(a_6), 0)$ is unstable and the predator population is persistent. Assume $g_6(P_0)_6 < b$.

Then, $Q_6(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_6 + u_a = -(\mu_6(a_6) + (f_{60})_6)u_6 - (\alpha_6)_6(a_6)N_6$$

$$= -(\mu_6(a_6) + f_{60})u_6 - \alpha_6(a_6)\int_0^\infty u_6(t, a_6)da_6$$

$$u_6(t, 0) = \int_0^\infty \beta_6\,(a_6)u_6(t, a_6)da_6$$

Using the method developed by Cushing [2, 3], we suppose

$$u_6(t, a_6) = \tilde{\mathrm{u}}_6(a_6)^{e^{\gamma(t-a_6)}}$$

And define $w = \int_0^\infty \tilde{\mathrm{u}}(a)e^{-\gamma a_6}\, da_6$. Then,

A solution of the first differential equation for $\tilde{\mathrm{u}}(a_6)$ gives,

$$\tilde{\mathrm{u}}_6(a_6) = e^{-\int_0^a \mu_6(\tau_6)d\tau - (f_{60})_6 a}\left(\tilde{\mathrm{u}}_6(0) - w\int_0^a e^{\int_0^\tau \mu_6(s_6)ds + (f_{60})}\tau_6\,(\alpha_6)_6(\tau_6)e^{\gamma_6\tau}d\tau\right).$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_6(a_6)e^{-\gamma_6 a_6 - (f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6)d\tau}\right)\tilde{\mathrm{u}}_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a(\alpha_6)_6(\tau_6)e^{-\gamma_6\tau_6}e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds}d\tau_6\,da_6$$

$$=0 \qquad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}da\right)(\tilde{u}_6(0))_6$$

$$+\left(-6-\int_0^{\infty}e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(\tau_6)d\tau}\int_0^a e^{(f_{60})_6\tau_6+\int_0^\tau \mu_6(s_6)ds+y\tau}\alpha_6(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_6 a_6-(f_{60})_6 a_6-\int_0^a \mu_6(t_6)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_6\,(a_6)Y_6(\gamma_6,a_6)da_6\right)u_6(0)$$

$$+w\int_0^{\infty}\beta_6(a_6)Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6=0,$$

(24)

$$\left(\int_0^{\infty}Y_6(\gamma_6,a_6)da_6\right)\tilde{u}_6(0)$$

$$+\left(-1-\int_0^{\infty}Y_6(\gamma_6,a_6)\int_0^a(\alpha_6)_6(\tau)_6(Y^{-6})_6(\gamma_6,\tau_6)d\tau_6\,da_6\right)w=0.$$

Equqtion (20) has a solution of the form of $u_6(t,a_6)=u_6(a_6)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_6(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) Y^{-6}(\gamma_6, a_6) d\tau_6$. We obtain the characteristic equation

$$f_{60}^{l}{}_6 \left( \int_0^\infty Y_6(\gamma_6, a_6) da_6 \right) \left( \int_0^\infty \beta_6(a_6) Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$+ \left( 1 - \int_0^\infty \beta_6(a_6) Y_6(\gamma_6, a_6) da \right) \left( 1 + (f_{60}^{l})_6 \int_0^\infty Y_6(\gamma_6, a_6) I_6(\gamma_6, a_6) da_6 \right)$$

$$= 0. \qquad (25)$$

Since

$$(\rho_0)_6(a_6) = (\rho_0)_6(0) e^{-\int_0^a \mu_6(\tau_6) d\tau_6 - (f_{60})_6 a_6}$$

$$I_6(\gamma_6, a_6) = \int_0^a (\rho_0)_6(\tau_6) e^{\gamma_6 \tau_6 + (f_{60})_6 \tau + \int_0^\tau \mu_6(s_6) ds_6} d\tau$$

$$= \frac{(\rho_0)_6(0)}{\gamma_6} (e^{\gamma_6 a_6} - 1)$$

Let

$$E_6(a_6) = e^{-(f_{60})_6 a_6 - \int_0^a \mu_6(\tau_6) d\tau_6}$$

This relationship is cardinal one. It is useful in quantum computation and other black hole dynamics problems

**Quantum space and quantum time**

(3758) **Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.**

(3759) **In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.**

(3760) **They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.**

(3761) **When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.**

(3762) **Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.**

**The Planck scale**

(3763) **Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".**

(3764) **Here ħ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.**

(3765) **The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.**

(3766) **Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)**

(3767) **Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.**

## Physical applications

**Some phenomena where quantum gravity plays a major role are the following.**

**The microscopic structure of spacetime**

(3768) **If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.**

(3769) **Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length**

(3770) **Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.**

(3771) **A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.**

**Early cosmology**

(3772) **According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.**

(3773) **However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.**

(3774) **A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.**

(3775) **Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".**

(3776) **One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.**

**Black holes**

(3777) **Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.**

(3778) **First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.**

(3779) **Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.**

(3780) **What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.**

(3781) **Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.**

**Astrophysical effects**

(3782) **Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.**

(3783) **An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.**

(3784) **The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.**

## Main approaches

(3785) **The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.**

(3786) **Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.**

(3787) **The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.**

## String theory

(3788) **String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.**

(3789) **Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.**

(3790) **String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.**

(3791) **Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.**

(3792) **However, when summed the entire series appears to be (=) divergent.**

(3793) **In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.**

(3794) **Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.**

(3795) **Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.**

(3796) **In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.**

## Loop quantum gravity and spin foams

(3797) **Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.**

(3798) **The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity.**

(3799) **Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime.**

(3800) **The granular properties of space can be (=) explicitly computed.**

(3801) **In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory.**

(3802) **The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks.**

(3803) **Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure.**

(3804) **The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit.**

(3805) **Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism.**

(3806) **This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism.**

## Noncommutative geometry

(3807) **Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime.**

(3808) **Quantum mechanics teaches us that dynamical quantities are (=) noncommutative -- meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed.**

(3809) **It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry.**

(3810) **A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces.**

(3811) **Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces.**

(3812) **This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity.**

**Others**

(3813) **Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays**

(3814) **. Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection.**

(3815) **Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible.**

(3816) **The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.**

**Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016**

(3817) **These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field.**

(3818) **The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets.**

(3819) **Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state.**

(3820) **Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands.**

(3821) **The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI**

(3822) **The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9].**

(3823) **This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive.**

(3824) **More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.**

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

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| (3825) | **Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.** |
| (3826) | **Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.** |
| (3827) | **The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.** |
| (3828) | **This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.** |
| (3829) | **The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.** |
| (3830) | **This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.** |
| (3831) | **For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.** |
| (3832) | **The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society** |

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

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| (3833) | **Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.** |
| (3834) | **We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.** |
| (3835) | **When this is true—as is often the case—the system serves as (=) a "bath" for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called thermalization.** |
| (3836) | **However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.** |
| (3837) | **Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1).** |
| (3838) | **Bringing them to thermal equilibrium requires (e) an external bath.** |

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| (3839) **But so far, there has been no rigorous proof of (e) the existence of such many-body localization. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption.** |
| (3840) **The concept of localization was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations).** |
| (3841) **The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century.** |
| (3842) **But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information.** |
| (3843) **Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment.** |
| (3844) **For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood.** |
| (3845) **It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7].** |
| (3846) **Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature.** |
| (3847) **One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.** |
| (3848) **However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].** |
| (3849) **Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.** |
| (3850) **Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.** |
| (3851) **As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked.** |
| (3852) **His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice.** |
| (3853) **An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site.** |
| (3854) **This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself).** |

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| (3855) **This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition.** |
| (3856) **However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime:** |
| (3857) **These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects.** |
| (3858) **It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies).** |
| (3859) **This is a realistic scenario for (e) all known systems.** |
| (3860) **On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects.** |
| (3861) **On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other.** |
| (3862) **No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true.** |
| (3863) **His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found.** |
| (3864) **Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization?** |
| (3865) **Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects?** |
| (3866) **Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions.** |
| (3867) **So in those systems, the quantum phase transition does indeed exist.** |
| (3868) **But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension.** |
| (3869) **Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters.** |

**Focus: Biological Cells Form Electric Circuits July 1, 2016• Physics 9, 75 Cells that are electrically active and that also produce light for easy voltage monitoring could lead to new studies of heart arrhythmias and possibly bio-computing. Figure caption expand figure H. M. McNamara et al., Phys. Rev. X (2016) Cell circuit. Waves of electrical activity propagate in a loop within a collection of cells engineered with electrical and optical properties (see video below).Some sentences, pictures and videos are deleted due to constraints on space.**

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| (3870) **The human heartbeat is produced by (e) electrical pulses that propagate through (e&eb) cardiac tissue, causing (eb) rhythmic muscle contraction.** |
| (3871) **Researchers have previously engineered cells to create (eb) an artificial tissue capable of producing (eb) coordinated electrical activity, and now a team has added the ability to monitor (e&eb) their electrical state by (e) detecting fluorescent emission.** |
| (3872) **They have also fashioned the cells into (e&eb) "living circuits" that might act as (=) model systems for studying heart behavior.** |
| (3873) **Constructing electrically active tissues from their basic building blocks by (e) genetically engineering cells allows (eb) researchers to study a "clean" system similar to (e) the more complicated cardiac tissue.** |
| (3874) **These systems let them investigate, for example, the electrical signals that lead to (eb) the potentially fatal irregular heartbeat called arrhythmia.** |
| (3875) **Previously researchers have introduced (eb) ion channels—proteins that conduct ions through (e&eb) cell membranes—into human kidney cells to make (eb) electrically "excitable" tissues [1].** |
| (3876) **Adam E. Cohen of Harvard University and his co-workers decided to add more ingredients to such cells, so that the electrical activity could be (=) light-activated and easily visualized.** |
| (3877) **They incorporated four additional gene products into (e&eb) human embryonic kidney (HEK) cells.** |
| (3878) **Two were ordinary ion channels; the other two were (=) a blue-light-activated ion channel (a light-sensitive switch) and (e&eb) a protein that fluoresces with (e&eb) an intensity that depends on (=) the voltage across the membrane in which it is embedded.** |
| (3879) **Light activation of electrically excitable cells has been shown before [2]. But Cohen and colleagues have now included a convenient "readout" mechanism (the fluorescent protein) and have used patterning methods to shape (e&eb) the cell cultures into "tissue circuits."** |
| (3880) **They grew rings of cells a few millimeters in diameter that could act as (e&eb) oscillators that support (eb) circulating waves of electrical excitation.** |
| (3881) **These waves arise from (e) the collective behavior of the cells: a pulse produced by (e) one cell can trigger another from (e) its neighbors, eventually producing (eb) a coherent wave of excitation traveling through (e&eb) the tissue.** |
| (3882) **It's this coordination, in a repeating cycle of excitation and quiescence, which gives rise to (eb) rhythmic activity in cardiac tissue.** |

**H. M. McNamara et al., Phys. Rev. X (2016)**

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| (3883) **A pattern of blue light (not visible) illuminates these cells, defining (eb) a rectangular loop around which electrical waves can propagate.** |
| (3884) **At about 0:07 the rectangle begins to shrink in (eb) the vertical direction, causing t(eb) he cycle time to gradually decrease. (In other experiments, the cells were grown in (eb) fixed, ring-shaped patterns.)** |
| (3885) **The ring of engineered HEK cells acted like an electronic "ring oscillator," a device with a characteristic frequency to (e&eb) its circulating electrical signal.** |
| (3886) **The electrical waves in a ring oscillator can propagate (e&eb) in either direction around the ring, enabling (eb) it to serve as a binary memory unit.** |

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| (3887) **Cohen and colleagues found that in their HEK tissues, these circulating waves remained (eb) stable for more than two hours, with (e&eb) the pulses completing a few circuits every second.** |
| (3888) **The electrical signals were easily visualized as (=) waves of fluorescent emission.** |
| (3889) **The researchers say that these tissues could be used to (e) understand and predict the behavior of real cardiac cells, especially arrhythmias.** |
| (3890) **"In real heart, if you pace the tissue fast enough, *(eb) it starts skipping beats and eventually becomes (=) arrhythmic," says Cohen.** |
| (3891) **This state is a well-known precursor to (e) sudden cardiac death and is associated with (e&eb) the breakup of the usual regular waves of electrical excitation.** |
| (3892) **Some drugs also have the dangerous side effect of inducing (eb) arrhythmia, and Cohen says that the new system might offer a quick and convenient way of screening for such hazards. The work "is really very elegant and deserves attention," says physiologist Gil Bub of Oxford University in the UK. "It opens new research questions for people like me who are interested in wave dynamics in excitable systems."** |
| (3893) **There might also be practical uses for information processing in living tissues. "By combining patterns of cell growth and (e&eb) patterns of ion channel expression, we think we can build diodes, amplifiers, and other simple logic elements," says Cohen.** |
| (3894) **Andrew Adamatzky of the University of the West of England in Bristol, a specialist in "unconventional computing," foresees circuits that perform computations through (e&eb) the interactions of electrical waves. "This research might lead to fabrication of implantable bio-computers," he says. This research is published in Physical Review X. –Philip Ball Philip Ball is a freelance science writer in London and author of Invisible: The Dangerous Allure of the Unseen (University of Chicago Press, 2014).** |

**Optically Controlled Oscillators in an Engineered Bioelectric Tissue Harold M. McNamara, Hongkang Zhang, Christopher A. Werley, and Adam E. Cohen Phys. Rev. X 6, 031001 – Published 1 July 2016**

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| (3895) **Complex electrical dynamics in excitable tissues occur throughout (e&eb) biology, but the roles of individual ion channels can be difficult to determine due to (e) the complex nonlinear interactions in (eb) native tissue.** |
| (3896) **Here, authors ask whether they can engineer a tissue capable of (e) basic information storage and processing, where (e) all functional components are known and well understood.** |
| (3897) **Authors develop a cell line with (e&eb) four transgenic components: two to enable (eb) collective propagation of electrical waves and (e&eb) two to enable optical perturbation and (e&eb) optical readout of membrane potential.** |
| (3898) **They pattern the cell growth to define (eb) simple cellular ring oscillators that run stably for >2 h (∼104 cycles) and that can store data encoded in (eb) the direction of electrical circulation.** |
| (3899) **Using (e) patterned optogenetic stimulation authors probe the biophysical attributes of (e) this synthetic excitable tissue in detail, including (e) dispersion relations, (e&eb) curvature-dependent wave front propagation, (e&eb) electrotonic coupling, and (e&eb) boundary effects.** |
| (3900) **Authors then apply the biophysical characterization to (e) develop an optically reconfigurable bioelectric oscillator.** |
| (3901) **These results demonstrate the feasibility of engineering bioelectric tissues capable of (e&eb) complex information processing with optical input and output. Received 12 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevX.6.031001 This article is available under the terms of the** |

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**Scientists detect radio emission from a nearby brown dwarf July 6, 2016 by Tomasz Nowakowski report**

(3902) **K2 and Spitzer images of W0607+24: The position of W0607+24 is marked by (e) a red circle. SDSS J060738.06+242939.5, marked by (e) an orange circle, are the brightest star in the K2 image.**

(3903) **The three faint (r = 18.3 − 19.6) stars marked by cyan circles, contribute significant flux to (e) the K2 measurement but are (-) negligible at Spitzer's longer wavelength and (e &êb) higher resolution. Other stars with r < 20 in SDSS are marked with blue circles. Credit: Gizis et al., 2016.**

(3904) **(Phys.org)—Located some 23 light years away, the brown dwarf designated WISEP J060738.65+242953.4 or W0607+24 for short, turns out to be (=) a source of radio emission.**

(3905) **According to a research paper published July 4, this substellar object showcases (eb) quiescent radio emission, making it one of (e) the most radio-faint, ultra-cool dwarfs yet detected.**

(3906) **Brown dwarfs like W0607+24 are objects that are too large to be called planets and too small to be stars. With a mass below that necessary to maintain hydrogen-burning nuclear fusion reactions, they are much cooler and dimmer than (eb) main sequence stars.**

(3907) **These objects could be also sources of (e) radio emission, but the cause of this process is still not completely understood.**

(3908) **Due to its proximity, W0607+24 is (=) an excellent target for studying the nature of (e) brown dwarfs.**

(3909) **It is the nearest known late-L dwarf in the northern hemisphere, and the third-nearest in (=) the whole sky, (eb) allowing researchers to study phenomena occurring in brown dwarfs.**

(3910) **That's why a team of astronomers, led by John Gizis of the University of (e)Delaware, Newark, carried out a series of mid-infrared, radio, and optical observations of W0607+24 to study any activity in this curious substellar object.**

(3911) **Motivated by the results from NASA's prolonged Kepler mission, named K2, which obtained (eb) long time-series photometry of W0607+24; they observed this object simultaneously with the Karl G. Jansky Very Large Array (VLA) in New Mexico and NASA's Spitzer Space Telescope.**

(3912) **The most important results were provided by VLA, as the team found that W0607+24 is (=) a source of quiescent, non-bursting, radio emission.**

(3913) **The scientists noted that the discovery of this process in this nearby object could be (=) crucial for our understanding of this process in (eb) ultracool dwarfs in general.**

(3914) **"W0607+24 is a valuable test case for (e) understanding the mechanisms of radio emission in ultracool dwarfs.**

(3915) **Some of the radio-detected objects emit (eb) rapid, intense bursts with high circular polarization, often found to occur periodically and attributed to coherent emission due to (e) the electron cyclotron maser instability," the paper reads.**

(3916) **The astronomers also suggested that the relative radio faintness of W0607+24 may be due to (e) its orientation if (e) the emission is concentrated at the poles.**

(3917) **However, follow-up observations are needed to better characterize (eb) the detected radio emission and to eliminate (e) all the uncertainties.**

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| (3918) **The team also hopes for more significant findings when (e) studying this nearby substellar object in the future.** |
| (3919) **"As the nearest known northern hemisphere late-L dwarf, this offers many possibilities for future studies. For example, it could serve as (=) an useful standard for (e) line-broadening and (e&eb) polarization studies.** |
| (3920) **If there is a planetary system aligned with (e&eb) (the primary's rotation axis, then we would expect very little radial velocity signal but a strong astrometric signal," the researchers concluded.** |
| (3921) **Moreover, the team plans to monitor (e&eb) additional late-M and L dwarfs in other K2 campaigns, expecting even more promising results.** |
| (3922) **They are more distant than W0607+24, but because they are warmer, (eb) some are brighter and higher signal-to-noise than that described in this paper.** |
| (3923) **In addition to detecting the radio emission, (eb)the scientists also estimated the age, mass and radius of W0607+24** |
| (3924) **. According to the study, the object is less than two billion years old, has (e) a maximum mass of about 0.055 solar masses and a radius approximately of 0.1 solar radii.** |
| (3925) **More information: WISEP J060738.65+242953.4: A Nearby. Pole-On L8 Brown Dwarf with Radio Emission, arXiv: 1607.00943 [astro-ph.SR] arxiv.org/abs/1607.00943** |
| (3926) **Authors present a simultaneous, multi-wavelength campaign targeting the nearby (7.2 pc) L8/L9 (optical/near-infrared) dwarf WISEP J060738.65+242953.4 in (eb) the mid-infrared, radio, and optical.** |
| (3927) **Spitzer Space Telescope observations show no (e) variability at the 0.2% level over 10 hours each in the 3.6 and 4.5 micron bands.** |
| (3928) **Kepler K2 monitoring over 36 days in Campaign 0 rules out (e) stable periodic signals in (eb) the optical with amplitudes great than 1.5% and periods between 1.5 hours and 2 days.** |
| (3929) **Non-simultaneous Gemini optical spectroscopy detects lithium, constraining (e) this L dwarf to be less than ~2 Gyr old, but no Balmer emission is observed.** |
| (3930) **The low measured projected rotation velocity (v sin i < 6 km/s) and lack of variability are (=) very unusual compared to other brown dwarfs, and authors argue that this substellar object is (=) likely viewed pole-on.** |
| (3931) **Authors detect quiescent (non-bursting) radio emission with (e&eb) the VLA.** |
| (3932) **Amongst radio detected L and T dwarfs, it has (e) the lowest observed L_nu and the lowest v sin i.** |
| (3933) **They discuss the implications of pole-on detection for (e) various proposed radio emission scenarios. © 2016 Phys.org** |

**A new look at the galaxy-shaping power of black holes July 6, 2016**

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| (3934) **It wasn't until the advent of X-ray astronomy that the full picture of (e) the galaxy-shaping power of black holes began to emerge (eb) with the ability to see plasma.** |
| (3935) **Data from a now-defunct X-ray satellite is providing (eb) new insights into the complex tug-of-war between galaxies, the hot plasma that surrounds them, and the giant black holes that lurk in their centres.** |

(3936) **Launched from Japan on February 17, 2016, the Japanese space agency (JAXA) Hitomi X-ray Observatory functioned for just over a month before contact was lost and (e&eb) the craft disintegrated.**

(3937) **But the data obtained during those few weeks was enough to paint (eb) a startling new picture of the dynamic forces at work within galaxies.**

(3938) **New research, published in the journal Nature today, reveals data that shows just how important the giant black holes in galactic centres are to the evolution (eb) of the galaxies as a whole.**

(3939) **"We think that supermassive black holes act(e&eb) like thermostats," said Brian McNamara, University Research Chair in Astrophysics at the University of Waterloo. "They regulate (e&eb) the growth of galaxies."**

**Champagne bubbles of plasma**

(3940) **During its brief life, the Hitomi satellite collected X-ray data from (e) the core of the Perseus cluster, an enormous gravitationally-bound grouping of (e) hundreds of galaxies.**

(3941) **Located some 240 million light years from earth, the Perseus cluster is one of (e) the largest known structures in the universe.**

(3942) **The cluster includes not only the ordinary matter that makes up the galaxies, but an (e) "atmosphere" of hot plasma with a temperature of tens of millions of degrees, as well as a halo of (e) invisible dark matter.**

(3943) **Earlier studies, going back to the 1960s, have shown that each of the galaxies in the cluster - and indeed most galaxies - likely contains (e) a supermassive black hole in its centre, an object 100 million to more than ten billion times as massive as our sun**

**New look at the galaxy-shaping power of black holes**

(3944) **Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving.**

(3945) **"These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science.**

(3946) **"Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy."**

(3947) **This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return.**

(3948) **The remaining matter gets absorbed into the black hole, adding to (e) its mass.**

(3949) **The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne.**

(3950) **The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution.**

(3951) **Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The**

**answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form.**

(3952) **"Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara.**

(3953) **"It creates (eb) these bubbles, and the bubbles keep the plasma hot.**

(3954) **That's what prevents galaxies from (e) becoming even bigger than they are now."**

(3955) **Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.**

(3956) **In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.**

(3957) **Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.**

(3958) **"It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.**

(3959) **"It's like nature's thermostat, which keeps these galaxies from (e) growing.**

(3960) **If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."**

**A new look at the galaxy-shaping power of black holes**

(3961) **An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the**

(3962) **McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.**

(3963) **Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.**

**Doomed satellite missions**

(3964) **Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.**

(3965) **At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.**

(3966) **Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a**

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| **leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.** |
| (3967) **The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.** |
| (3968) **These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.** |
| (3969) **"But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people** |
| (3970) **. It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."** |
| (3971) **The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.** |
| (3972) **"This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website** |

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

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| (3973) **Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.** |
| (3974) **They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.** |
| (3975) **However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.** |
| (3976) **Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.** |
| (3977) **X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well.** |
| (3978) **The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma.** |
| (3979) **These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6.** |
| (3980) **Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus.** |

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| (3981) | **A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate.** |
| (3982) | **Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure.** |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (3983) | **This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy.** |
| (3984) | **Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole.** |
| (3985) | **But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy.** |
| (3986) | **A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole.** |
| (3987) | **It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth.** |
| (3988) | **"There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA).** |
| (3989) | **This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies.** |
| (3990) | **However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated.** |
| (3991) | **In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1.** |
| (3992) | **We know dark matter exists only from (e) its gravitational effects.** |
| (3993) | **It holds together galaxies and (e&eb) galaxy clusters.** |
| (3994) | **Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years.** |
| (3995) | **To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies.** |
| (3996) | **They used star motions as tracers to (e) weigh the galaxies' central black holes.** |
| (3997) | **X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto.** |
| (3998) | **They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone.** |
| (3999) | **This connection is likely to be related to (e&eb) how elliptical galaxies grow.** |

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| (4000) **An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together.** |
| (4001) **Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole.** |
| (4002) **"In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp** |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (4003) **The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the …** |
| (4004) **Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness?** |
| (4005) **Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness.** |
| (4006) **"Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain."** |
| (4007) **Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness.** |
| (4008) **"With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication."** |
| (4009) **Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain** |

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| **changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories.** |
| (4010) **The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen.** |
| (4011) **In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer.** |
| (4012) **Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants.** |
| (4013) **These were treated as (=) "aware" and "unaware" trials, respectively.** |
| (4014) **Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it.** |
| (4015) **The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another.** |
| (4016) **Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion.** |
| (4017) **No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness.** |
| (4018) **"We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said.** |
| (4019) **"Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness."** |
| (4020) **The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner.** |
| (4021) **Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain.** |
| (4022) **"We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said.** |
| (4023) **"This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients** |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

(4024) **Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter.**

(4025) **As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined.**

(4026) **Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions.**

(4027) **"Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.**

(4028) **The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.**

(4029) **However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.**

(4030) **The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.**

(4031) **Noda and his colleagues instead studied the change in emissions over (e&eb) time.**

(4032) **By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.**

(4033) **Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).**

(4034) **The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.**

(4035) **Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.**

(4036) **"The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp**

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(4037) **The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking "palm tree" organization.**

(4038) **The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society**

(4039) **(Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible?**

(4040) **With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem.**

(4041) **This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other.**

(4042) **According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible.**

(4043) **This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges.**

(4044) **In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century.**

(4045) **Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching.**

(4046) **In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom.**

(4047) **Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth.**

**Spheres on trees**

(4048) **Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration.**

(4049) **They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement.**

(4050) **Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum.**

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| (4051) **By visualizing the problem in this way, the researchers noticed that these particular graphs don't (e) have lots of branches extending from lots of (e) other branches.** |
| (4052) **Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure.** |
| (4053) **The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network.** |
| (4054) **As a result, it provides (eb) an important clue for finding the global minimum.** |
| (4055) **It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum.** |
| (4056) **This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large.** |

**Thompson's problem 2**

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| (4057) **By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society** |
| (4058) **The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima.** |
| (4059) **The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node.** |

**Implications for Thomson's and Smale's problems**

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| (4060) **Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before.** |
| (4061) **"Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around."** |
| (4062) **As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem.** |
| (4063) **"I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view."** |
| (4064) **In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges.** |
| (4065) **From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear.** |

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| (4066) | **"While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK.** |
| (4067) | **"For larger numbers, there will likely be (=) multiple funnels.** |
| (4068) | **Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum."** |
| (4069) | **Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.** |
| (4070) | **"The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.** |
| (4071) | **"Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org** |

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**Introduction:**

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price, University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss.**

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| **(673)** | **Black Holes are regions of (e) space in which gravitational fields are so strong that no particle or signal can escape (eb) the pull of gravity.** |
| **(674)** | **The boundary of this no-escape region is (=) called the event horizon, since distant observers outside the black hole cannot (e) see (cannot get light from) events inside.** |
| **(675)** | **Although the fundamental possibility of such an object exists within (eb) Newton's classical theory of gravitation, Einstein's theory of gravity makes (eb)black holes inevitable under some circumstances.** |
| **(676)** | **Prior to the early 1960s, black holes seemed to be (=) only an interesting theoretical concept with no (e) astrophysical plausibility, but with (e&eb) the discovery of quasars in 1963 it became clear that very exotic astrophysical objects could exist.** |
| **(677)** | **Nowadays it is taken for granted that black holes do exist in (eb) at least two different forms.** |
| **(678)** | **Stellar mass black holes are (=) the endpoint of the death of some stars, and supermassive black holes are the result of (e) coalescences in the centers of most galaxies, including our own.** |
| **(679)** | **No signal can propagate from (e) inside a black hole, but the gravitational influence of a black hole is (=) always present.** |

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| (680) | This influence does not (e) propagate out of the hole; it is permanently present outside, and depends only on (=) the total amount of mass, angular momentum, and electric charge that have gone into forming (eb) the hole. |
| (681) | Black holes can be detected through the influence of this (e&eb) strong gravity on the surroundings just outside the hole. |
| (682) | In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them. |
| (683) | Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors. |

**Classical vs. relativistic black holes**

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| (684) | Something like a black hole exists within (eb) Newton's classical theory of gravity. |
| (685) | In that theory, an energy argument tells us that there is (=) an escape velocity $v_{esc}=\sqrt{2GM/R}$ from (e) the surface of any spherical object of (e) mass M and radius R |
| (686) | If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity. |
| (687) | Thus the condition for such an "unseeable" object is $R<2GM/c^2$ |
| (688) | In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape. |
| (689) | This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity. |
| (690) | To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it. |
| (691) | The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c . |
| (692) | This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed. |
| (693) | An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity. |
| (694) | Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable. |

**Stationary black holes**

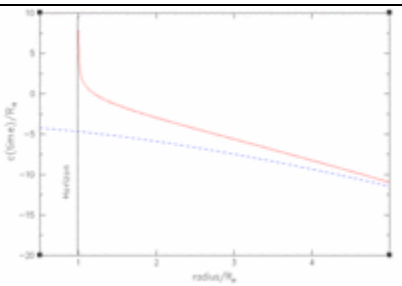
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| (695) | In the Newtonian theory, gravity is described by (e) the potential $\Phi$. |
| (696) | Inside a spherical object the form of (e) $\Phi(r)$ depends on (=) the interior structure, but in the vacuum outside (e) matter the potential $-GM/r$ depends only on (=) the interior mass. |
| (697) | Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object. |

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| **(698) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward.** |
| **(699) This horizon radius RH=2GM/c2 is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory.** |
| **(700) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution.** |
| **(701) Radius RHhere actually means that the area of the event horizon is (=) 4πR2H.** |
| **(702) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution.** |
| **(703) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity.** |
| **(704) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own.** |
| **(705) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole.** |
| **(706) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density.** |
| **(707) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed.** |
| **(708) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole.** |
| **(709) In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter.** |
| **(710) Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum.** |
| **(711) In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior.** |
| **(712) The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object.** |
| **(713) Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum.** |
| **(714) These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.]** |
| **(715) All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time.** |
| **(716) But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?"** |
| **(717) The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point.** |

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| **(718) One of these coordinates is (=) called the "coordinate time."** |
| **(719) Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate.** |



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| **(720) Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole** |
| **(721) At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole.** |
| **(722) Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations.** |
| **(723) For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree** |
| **(724) An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events.** |
| **(725) The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving.** |
| **(726) Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle.** |
| **(727) The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return.** |
| **(728) But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon.** |
| **(729) Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future.** |
| **(730) The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?"** |

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| **Black hole parameters** |

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| **(731) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin).** |
| **(732) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun.** |
| **(733) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits.** |
| **(734) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form.** |
| **(735) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds.** |
| **(736) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2.** |
| **(737) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed.** |
| **(738) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer.** |
| **(739) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2)** |
| **(740) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials.** |
| **(741) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable** |
| **(742) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite.** |
| **(743) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole.** |
| **(744) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars.** |
| **(745) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity.** |
| **(746) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity.** |
| **(747) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite.** |
| **(748) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity.** |
| **(749) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum.** |

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| **(750) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship.** |
| **(751) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= (4πG2/c4)[(M+M2−a2−−−−−−−√)2+a2].** |
| **(752) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored.** |

**Dynamical black holes**

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| **(753) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape.** |
| **(754) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, a horizon can be born.** |
| **(755) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid.** |
| **(756) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward.** |
| **(757) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations.** |
| **(758) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes.** |
| **(759) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3.** |
| **(760) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/** |
| **(761) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers.** |
| **(762) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission.** |
| **(763) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.** |
| **(764) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.** |
| **(765) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.** |
| **(766) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.** |

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| **(767) The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature 6×10−8 (Mo/M) K.** |
| **(768) Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.** |

Then $[Y_6]_6[\gamma_6]_6, a_6) = E_6(a_6)e^{-\gamma_6 a_6}$. by a simple calculation,

$$\int_0^\infty \beta_6(a_6)Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma}\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}\ da\right);$$

$$\int_0^\infty Y_6(\gamma_6, a_6)I(\gamma_6, a_6)da = \frac{[\rho_0]_6(0)}{\gamma_6}\left(\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}\ da_6\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_6(0)}{\gamma_6}[f_{60}^6]\int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}da_6\ (\int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6)$$

$$+\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right).$$

$$\left(1 + \frac{(\rho_0)_6(0)}{\gamma_6}(f_{60}^{'})_6\left(\int_0^\infty E_6(a_6)da_6 - \int_0^\infty E_6(a_6)e^{-\gamma_6 a_6}da_6\right)\right) = 0.$$

Hence the characteristic equation becomes

$$\left(1 - \int_0^\infty \beta_6(a_6)E_6(a_6)e^{-\gamma_6 a_6}da_6\right)(1 + \frac{(f_{60}^{'})_6}{\gamma_6}(\rho_0)_6(0)\int_0^\infty E_6(a_6)da_6) = 0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1 - \int_0^\infty \beta_6(a_6) Y_6(\gamma_6 a_6) da_6 = 1 - \int_0^\infty \beta_6(a_6) E_6(a_6) e^{-\gamma_6 a_6} da_6 \neq 0,$$

Since $\int_0^\infty \beta_6(a_6) Y_6(\gamma_6 a_6) \int_0^a \alpha_6(\tau_6) Y^{-6}(\gamma_6, \tau_6) d\tau_6 da_6 > 0$. Therefore, the characteristic solution is

$$\gamma_6 = -(f_{60}^{'})_6 (\rho_0)_6 \int_0^\infty E_6(a_6) da_6.$$

Since $f_6'(P_0)_6 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_6), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE OF DEBIT TYPES AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_6), 0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_6 + (\mu_6(a_6) + [f_6]_6(P_6^*) + [f_2]_6(Q_6^*))P_6^* = 0, \quad (26)$$

$$p_6^*(0) = h_6(Q_6^*) \int_0^\infty \beta_6(a_6) p_6^*(a_J) p_6^*(a_6) da_6 \quad (27)$$

$g_6(P_6^*) - b_6 = 0,$ (28)

Where $P_6^*$=$\int_0^\infty p_6^*(a_6)\, da_6$

Equation (28) yields $P^* = g^{-6}(b)$,

From (26),

$p_6^*(a) = p_6^*(0)e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6}$

Substituting into (27), we need to solve

$$\left(1 - h_6(Q_6^*)\int_0^\infty \beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(Q_6^*)a_6} da_6\right) p_6^*(0) = 0$$

For $Q_6^*$

Fix $p_6^*(0)$ and $P_6^*$ Let

$$H_6(z_6) = 1 - h_6(z_6)\int_0^\infty \beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6} da_{,6}$$

And assume that

$$[k_3]_6 \int_0^\infty \beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(g_6^{-6}(b_6))a_6 - [k_2]_6 a_6}\, da_6 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_6(z_6)] > 0.$$

Clearly,

$$\frac{dH_6}{dz_6} = h_6(z)([f_2^{'}]_6(z_6))\int_0^{\infty} a_6\beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - f_6(P_6^*)a_6 - f_2(z_6)a_6}\, da_6$$

$-h^{'}(z_6)\int_0^{\infty}\beta_6(a_6)\, e^{\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6 - [f_2]_6(z_6)a_6}\quad da_6$

By the assumptions on $\left[f_2^6\right]_6 and\ [h^6]^6 dH_6/dz > 0$. Hence, if

$$H_6(0) = 1 - \int_0^{\infty}\beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6(P_6^*)a_6}da_6 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_6(0) < 0$ follows if $g_6^{-}(b_6) < [f_6^{-6}]_6(r_6)$, where $r_6$ is the solution of (12).

Indeed, since

$$R_6(s_6) = \int_0^{\infty}\beta_6(a_6)\, e^{-\int_0^a \mu_6(\tau_6)d\tau_6 e^{-y_6^a}}da_6$$

Is a decreasing function of $s_6$ when $P_6^* = [g^{-6}]_6(b_6) < [f_6^{-6}]_6(r_6)$, that is $[f_6]_6(P_6^*) < r$, $R\left(f_6(P^*)\right) > 1$which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_6(g^{-6}(b))) > 6, i.e.,$

$$\int_0^\infty \beta_6(a_6) e^{-\int_0^a \mu_6(\tau_6)d\tau_6 - [f_6]_6\left(g_6^{-6}(b)\right)a_6} > 1. \quad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_6^*(a_6), Q_6^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_6^*(a_6), Q_6^*)$ is more complicated. We again consider perturbations $m_6(t_6 a_6) = p_6(t_6, a_6) - p_6^*(a_6)$ and $W_6(t) = Q_6(t) - Q^*$.Then

$$M_6(t) \equiv \int_0^\infty m_6(t, a_6) da_6 = P_6(t) - P_6^*,$$

$$[m_6]_6 + [m_a]_6 = [p_6]_6 + [p_a]_6 - \left[p^*\right]_6(a_6) \approx -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6$$

$$-[f_{6*}^6]_6 p_6^*(a_6)M - [f_{2*}^6]_6 p_6^*(\mathrm{a})W_6,$$

$$m_6(t,0)p_6(t,0) - p_6^*(0) \approx h_* \int_0^\infty \beta_6(a_6) m_6(t, a_6)\, da + \frac{[h_*^6]_6}{[h_*]_6}\, p_6^*(0) W_6$$

$$\frac{dW_6}{dt} = \frac{dQ_6}{dt} \approx [g_*^6]_6 Q_6^* M_{,6}$$

Where $[f_{6*}]_6$ denoting $[f_6(P^*)]_6$, $[f_{6*}]_6$ denoting $[f_2(P^*)]_6$, $\left[f_{6*}^{'}\right]_6$ denoting $\left[f^{'}{}_6(P^*)\right]_6 \left[f_{2*}^{'}\right]_6$ denoting $[f_2]_6$ $(Q_6^*,$ $h_{*}^{}{}_6$ denoting $h_*^{'}$ denoting $h_6^{'}$ $(Q_6^*)$, and $\left[g_*^{'}\right]_6$

Denoting $\left[g^{'}\right]_6(P_6^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_6]_6 + [m_a]_6 = -(\mu_6(a_6) + [f_{6*}]_6 + [f_{2*}]_6)m_6 - [f'_{6*}]_6[p_6^*](a_6)M_6 - [f'_{2*}]_6 p_6^*(a_6)W_6$$

$$m_6(t,0) = [h_*]_6 \int_0^\infty \beta_6(a_6)m_6(t,a_6)da_6 + \frac{[h_*^6]_6}{[h^*]_6}[p^*]_6(0)\,W_{,6} \qquad (30)$$

$$\frac{dW_6}{dt} = [g_*^6]_6[Q_6^*]M_6$$

Integrating $[m_6]_6(t,a_6)$ with respect to a from 0 to ∞, we have

$$\frac{dM_6(t)}{dt} = \int_0^\infty [mt]_6(t,a_6)da_6$$

$$= -\int_0^\infty [m_a]_6\,(t,a_6)da_6 - \int_0^\infty \mu(a_6)m(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + [f'_{6*}]_6 \int_0^\infty p_6^*(a_6)da_6\right)M_6 - [f_{2*}^6]_6 \int_0^\infty p_6^*(a_6)da_6 W_6$$

$$= m_6(t,0) - \int_0^\infty \mu_6(a_6)\, m_6(t,a_6)da_6$$

$$-\left([f_{6*}]_6 + [f_{2*}]_6 + [f_{6*}']_6 P_6^*\right)M_6 - [f_{2*}']_6 P_6^* W_{,6}$$

From (30), we arrive at the equations

$$\frac{dW_6(t)}{dt} = \int_0^\infty ([h_*]_6\beta_6(a_6) - \mu_6(a_6))m_6(t,a_6)da_6$$

$$-\left([f_{6*}] + [f_{2*}]_6 + [f_{6*}']_6 P_6^*\right)M_6 + \left(\frac{[h_*']_6}{[h_*]^6}\,P_6^*(0) - f_{2*}' P_6^*\right)W_6,$$

(36)

$$\frac{dW_6(t)}{dt}[g^{'}{}_{*}]_6 Q_6^* M_{.6}$$

By the mean value theorem,

$$\int_0^{\infty}([h_{*}]_6\beta_6(a_6) - \mu_6(a_6))m_6(t,a_6)da_6 = ([h_*]_6\beta_6(\theta_6) - \mu_6(\theta_6))M_{,6}$$

Where $0< \theta_6 = \theta_6(t) < \infty$.

And if $a_6^2[a_6^2]_6<4\,[a_2]_6<4\,a_2 e$,

Re $\lambda_6[\lambda_6]_6 = \text{Re}\lambda_2[\lambda_2]_6 = -\frac{6}{2}a_6[a2\ ]_6 < 0$.

Let $\Delta=\max_{2<u<\infty}|h_*[h_*]_6\beta_6(a_6) - \mu_6(a_6)|$. Then$|n_6(t)|$ ⊠⊿∀t ⊠0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium($\rho^*(a), \boldsymbol{Q}^*$) is asymptotically stable if either $a_6^2 <4\,a_2 e$ and $[a_6] > 2\Delta$,or $[a_6^2]$ ⊠4$[a_2]$e and $(a_6 - \sqrt{(a_6^2)_6 - 4(a_2)_6 e}) > 2\Delta$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+6}\right)$ and $\beta_6(a_6) = ae^{-ma}$ Then $\Delta=[\mu_0] + 1$ Simple calculation leads to

$$\int_0^{\infty}\beta_6(a_6)e^{-\int_0^a \mu_6(\tau_6) + d\tau_6}da = \frac{6}{(m+\mu_0+1)^2} + \frac{2}{(m+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H1), we assume $m_6 + [\mu_0]$ 0.5. Assume$[f_6](x_6)= K_6x_6$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_6(x_6)= [0x_6]/(1 + x_6)$; the functional responses $[[f_2](x) = [0]\, x_6/(1 + x_6)$and $h_6(x_6)$=(1+$\omega_6 x_6$)/$(1 + x_6)$(See Holling [7] and Rashevsky [13])

Equation (64),now has the form of

$$\frac{m_6 + [\mu_0] + [\mathrm{K}][P_0])}{[m_6]_{\,6} + (\mu_0)_{\,6} + [\mathrm{K}][P_0])\,{}^{3}_{6}\ [\mathrm{K}]_{\,6}[P_0]\,{}^{3}_{6}} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(6 + P_0) < b_6$i e $[P_0] < b/(\delta_6 - b_6)$then the equilibrium$(P_0)$ $[P_0]_{\,6}(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_6$

$[f_{6*}] = K_6\ \mathrm{P}^{*}\ [f_{2*}] = \frac{\sigma \mathrm{Q}_6^*}{6+\mathrm{Q}_6^*}$ $[h_*] = \frac{6+\omega_6 \mathrm{Q}_6^*}{6+\mathrm{Q}_6^*}$

$[f_{6*}] = K_6\ \left[\mathrm{f}\,{}^{6}_{2*}\right] = \frac{\sigma_6}{(6+\mathrm{Q}_6^*)^2}[\ \mathrm{g}\,{}^{6}_{*}] = \frac{\sigma_6}{(6+\mathrm{P}_6^*)^2}$ $\left[\mathrm{h}\,{}^{6}_{*}\right] = \frac{\omega_6 - 6}{(6+\mathrm{Q}_6^*)^2}$

In order to have (H3) and (H4) hold,$m, \mu_0,$ K, b, δ, and $\omega$ need to satisfy

$$\frac{\omega_6}{(m + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 + b_6} + \sigma)^{\,2}} + \frac{2\omega_6}{([m_6]_{\,6} + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 + b_6} + \sigma_6)^{\,3}} < 1$$

And

$$\frac{1}{(m_6 + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 - b_6})^{\,2}} + \frac{2\omega_6}{(m_6 + [\mu_0]_{\,6} + 1 + \frac{K_6 b_6}{\delta_6 - b_6})^{\,3}} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_6 + [\mu_0]_6 + 1 + \frac{K_6 b_6}{\delta_6 - b_6} + (m_6 + [\mu_0]_6 + \frac{K_6 b_6}{\delta_6 - b_6} + \sigma_6) Q_6^*)^3$$

$$= m_6 + [\mu_0]_6 + 3 + \frac{K_6 b_6}{\delta_6 - b_6}(1 + Q^*)(1 + \omega_6 Q_6^*$$

Now

$$\left[a_6^2\right] - 4\left[a_2\right]e = \left(\frac{2K_6 b_6}{\delta_6 - b_6}\right)^2 + \frac{4K_6 b_6 \sigma_6 Q_6^*}{(\delta_6 - b_6)(1 + Q_6^*)} + \frac{\sigma_6^2 Q_6^{*2}}{(1 + Q_6^*)^2}$$

$$\frac{4 b_6 \sigma_6 (\delta_6 - b) Q^*}{\delta_6 (1 + Q_6^*)^2} - \frac{4 (\delta_6 - b_6)^2 (1 - \omega_6) Q_6^*}{\delta_6 (1 + \omega_6 Q_6^*) 1 + Q_6^*)} P * (0)$$

And

$$\Delta = \lim_{0 < a < \infty} |[h_*]_6 a e_6^{mu} [\mu_0]_6 - \frac{a_6}{a_6 + 6}$$

Hence if $[a_6^2 > 4(a_2)_6 e$ and

$$[a_6] = \frac{2K_6 b_6}{\delta_6 - b_6} + \frac{\sigma_6 Q_6^*}{6 + Q_6^*} > 2\Delta$$

The nontrivial equilibrium is stable.

IF we define

$$[a_6]_6 = [f_{6*}]_6 + [f_{2*}]_6 + [f_{6*}]\,P^*_6$$

$$[a_2]_6 = \left[f^6_{2*}\right]P^*_6 - \frac{h^6_*}{h_*}P^*(0)$$

$$e=[g^6_*]_6\,Q^*_6$$

$$\eta_6(t) = \left[h_*\right]_6\beta_6\left(\theta_6(t)\right) - \mu_6\left(\theta(t)\right),$$

The system(36)can be written as the following system:

$$[M(t)]\left[a_6\right]M_6 - [a_2]_6 W_6 - [n(t)]M_6$$

$$w_6(t) = eM_6$$

Write

$$A=\begin{pmatrix} -[a_6]_6 - [a_2]_6 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_6 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_6 \\ W_6 \end{pmatrix}$$

The solution of (32)is

$$X_6(t) = e^{A,t}\,[X]_6(0) + \int_0 e^{At}$$

Other problems to which the model is applied are: To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case β→∞ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness condition of (e) entropy. It is shown that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology

(gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure authors have discovered. It is shown that Legendre invariant thermodynamic scalar curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws, leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions,**

**geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged AdS black holes by considering the case of the RN-AdS black holes surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase transition. For the case $\omega_q=-2/3$ authors derive analytic expressions of (e) critical physical quantities, while (e) for cases $\omega_q\neq-2/3$ appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark

energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=)

chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross

section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions, while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center

collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma? Paul M. Chesler STOP Paul M. Chesler, Krishna Rajagopal** calculate how the energy and (e&eb) the opening angle of jets in N=4 SYM theory evolve as they propagate through (e&eb) the strongly coupled plasma of that theory. **Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb)

heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version) **Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb) strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion

collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar (=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI: 10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills

(SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the

decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed

black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to

equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a

nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about

the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic

map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstates thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11 August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical

advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.**

Characterizing the behaviour of strongly coupled quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at $\lambda=\lambda c$, the correlation length diverges as $\xi\sim|\lambda-\lambda c|-\nu$, while (e) the gap between the ground state and first excited state

vanishes as (e) ΔE∼ξ−z∼|λ−λc|zν. Exactly at the critical point λ=λc, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as ω∼|k−k0|z. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004 STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting

qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element. The basic architecture presented here can be expanded, offering (eb) the possibility for the coherent interaction of (e&eb) a large number of superconducting qubits. **Nature 449, 438-442 (27 September 2007) | doi:10.1038/nature06124; Received 18 April 2007; Accepted 25 July 2007 Coherent quantum state storage and transfer between two phase qubits via a resonant cavity Mika A. Sillanpää1, Jae I. Park1 & Raymond W. Simmonds1 STOP. Chui-Ping Yang, Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e)

long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007**

# PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION ATTRITION MODELS

## MOODULE NUMBER SEVEN

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Systems to which have applied the model are as follows: Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures have been proposed. An attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in**

**(eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP. In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011 http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 Qubit transfer**

between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys. Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer. Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP. In recent years S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P have investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys.

**Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP. Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2 [quant-ph] for this version) Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP. Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue from over (e&eb) a thousand plant and animal samples. NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP. Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this behavior, enhancing and accelerating (eb, eb+) the growth of correlations. Michael Kastner gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. Michael Kastner tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP. A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E.**

**Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91, 075435 – Published 26 February 2015 STOP. Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1 proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point'' with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. Membranes at quantum criticality Petr Hořava1, 2 Published 4**

**March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that (eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035 Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013 STOP Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study**

**the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen STOP Authors study the dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing**

**seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) Dynamic Service Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by H∗(X,(Λ∗TX)⊗g), where X is the target hyperk\"ahler manifold. We further give a mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the**

**low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition. Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb) the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537 © 2016 American Physical Society Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant (Δα/α) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and Δα/α to (e&eb) this conclusion.**

**However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts z≲3, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the Λ cold dark matter model value (jΛCDM=1); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in α(z), linked to (e&eb) a linear coupling (ζ) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order |ζ|∼10−4, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in Δα/α (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems, authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them**

**are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP Authors measure the star formation properties of (e) two large samples of galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the star formation rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a function of environment. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]. Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size √ n × √ n? The problem is that √n time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time O (√n) for d ≥ 3, or O (√nlog5/2n) for d = 2. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ∗This work was mostly done while the author was a PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship and by CIAR. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In**

**d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e) sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP The relationship between the foundations of mathematics and (e&eb) physics is a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb)**

**flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph] (or arXiv: 1503.05600v1 [math-ph] for this version) Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative**

**real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to the fine structure constant. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor goes to 0 as the (e&eb) time goes to 0, the big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph]**

**(or arXiv: 1306.4613v1 [quant-ph] for this version) Effects of gauge theory based number scaling on geometry Paul Benioff STOP In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x, outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph] (or arXiv: 1211.3381v2 [quant-ph] for this version) Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value, \int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances, the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an**

**observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include (e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP This work is based on (e) the field of reference frames based on (e) quantum representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc) Journal reference: Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul Benioff STOP A quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of**

**these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference: Journal of Physics: Conference Series 70 Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows: System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The**

**Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol**

## 68. MODEL FORMULATION

Let $p_7(t, a_7)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(g_7)Q_7$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_7(0) = 0, \quad g_7(P_7) > 0, \ P_7 \ 0, \qquad \lim_{P_7 \to \infty} g_7(P_7) = (K_7)_7 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_7 + (\rho_a)_7 + D_7(a_7, P_7, Q_7)\rho_7 + 0, \qquad 0 < a_7 \ \infty,$$

$$\rho_7(t,0) = \int_0^\infty n_7(a_7, P_7, Q_7)\rho_7(t_7, a_7)da_7$$

Where $D_7$the mortality or debit type function is function and $n_7$ is the fertility function.

Put $_7$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_7}{dt}\right]_7 = x_7 \left[\frac{dQ_7}{dt}\right]_7 - x_7 \left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_7}{dt}\right]_7 = x_3\left[\frac{dQ_3}{dt}\right]_7 - x_4\left[\frac{dQ_7}{dt}\right]_7$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_7}{dt}\right]_7 - x_6\left[\frac{dQ_3}{dt}\right]_7$$

Where $x_7, x_7, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_7$ and $g_7(P_7)Q_7$

Now when you add the term $(b_7Q_7)$ with $g_7(P_7)(Q_7)$ we are infect see is that $(b_7Q_7)$ is credited incrementally part of $g_7(P_7)(Q_7)$ or instant nearly, when subtracted from $\frac{dQ_7}{dt}$ , it means that $g_7(P_7)(Q_7)$ is slowly eaten by $\frac{dQ_7}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

Following equation throws light on important issues in the following cases

**Simulation suggests non-water based life could exist on Saturn's moon Titan July 5, 2016 by Bob Yirka report Titan True-color image of layers of haze in Titan's atmosphere. Credit: NASA**

(4072) (Phys.org)—A team of researchers at Cornell University has built and run a simulation that showed (eb) prebiotic reactions could possibly occur on (eb) the surface of one of Saturn's moons, Titan, suggesting (eb) the possibility of life evolving in a place where (e) it is too cold for water to be a factor.

(4073) In their paper published in Proceedings of the National Academy of Sciences, the team describes the simulation they created in response to (e&eb) the discovery (by the Huygens probe) that polymers such as polyimine might have already developed on (eb) the moon's surface.

(4074) As scientists continue to search for life on other planets, more and more they are beginning to concede that if it does exist, (eb) it may not necessarily exist in the Goldilocks, or habitual zone.

(4075) Such planets are all the "right" distance from (e) their star, and hopefully, also have water.

(4076) But recent evidence has suggested that some chemical reactions could lead to (eb) types of life forms where there is no water, which further suggests that it may exist outside of (e) what is now considered the habitual zone.

(4077) For life to come about in such places, researchers reason, there would likely need to be (=) some sort of action going on—and that is why there has been (=) so much focus on Titan; it is the only object in our solar system, besides Earth, that has both rainfall and erosion due to (e) liquid movement.

(4078) But the water it has is locked far underground and the moon is (=) too cold to support an impact by (e&eb) water anyway.

(4079) But, as the researchers with this new effort discovered after poring over data sent back by Huygens, the surface does have (e) hydrogen cyanide in its sediment, brought down from (e) the atmosphere by methane and ethane rain.

(4080) It was those molecules that drove (eb) the design of the simulations—the team wanted to see if they could form the basis of reactions that could lead to (eb) the creation of polymers such as polyimine, which the team notes are (=) conducive to the formation of prebiotic reactions leading to (eb) a form of life.

(4081) Their simulations showed (eb) that such reactions are possible and that the structures that came about were also able to absorb (e) sunlight in the wavelengths present on (eb) the Titan surface.

(4082) The researchers suggest their work, and that done by others indicates (eb) that a return to the planet by (e) a new probe might be in order, one able to test for (e) a different form of life, or at least its precursors. More information: Polymorphism and electronic structure of polyimine and its potential significance for prebiotic chemistry on Titan, PNAS, www.pnas.org/cgi/doi/10.1073/pnas.1606634113

**Abstract**

(4083) The chemistry of hydrogen cyanide (HCN) is believed to be central to (e) the origin of life question.

(4084) Contradictions between Cassini–Huygens mission measurements of (e) the atmosphere and the surface of (e) Saturn's moon Titan suggest (eb) that HCN-based polymers may have formed on the surface from (e) products of atmospheric chemistry.

(4085) This makes Titan a valuable "natural laboratory" for exploring potential nonterrestrial forms of (e) prebiotic chemistry.

(4086) Authors have used theoretical calculations to investigate the chain conformations of polyimine (pI), a polymer identified as (=) one major component of polymerized HCN in laboratory experiments.

(4087) Thanks to its flexible backbone, the polymer can exist in (eb) several different polymorphs, which are (=) relatively close in energy.

(4088) The electronic and (e&eb) structural variability among them is extraordinary.

(4089) The band gap changes over (e&eb) a 3-eV range when moving from a planar sheet-like structure to (e&eb) increasingly coiled conformations.

(4090) The primary photon absorption is predicted to occur in (eb) a window of relative transparency in Titan's atmosphere, indicating (eb) that pI could be photochemically active and (e&eb) drive chemistry on the surface.

(4091) The thermodynamics for adding and removing HCN from pI under Titan conditions suggests (eb) that such dynamics is plausible, provided that (e) catalysis or photochemistry is available to sufficiently lower reaction barriers.

(4092) Authors speculate that the directionality of pI's intermolecular and (e&eb) intramolecular =N–H…N hydrogen bonds may drive (eb) the formation of partially ordered structures, some of which may synergize with (e&eb)photon absorption and act catalytically. Future detailed studies on proposed

mechanisms and the solubility and density of the polymers will aid in the design of future missions to Titan. Journal reference: Proceedings of the National Academy of Sciences search and more info website © 2016 Phys.org Read more at: http://phys.org/news/2016-07-simulation-non-water-based-life-saturn.html#jCp

**Water waves travel back in time to retrace their ripples By Jacob Aron (New Scientist)**

(4093) Nudge your bath tub just right and (eb) you might end up reversing time.

(4094) Don't worry – you won't end up wet, naked and without a towel in the Cretaceous period. Emmanuel Fort at the Langevin Institute in Paris, France, and his colleagues have developed what they call a "time mirror", which can make water waves rewind into (e&eb) their past. Humans and rubber duckies are (=) affected.

(4095) To understand how this works, imagine dropping a stone in a pond and watching waves ripple out from a central point in an ever-expanding circle. The equations governing the behaviour of waves say that it should also be possible for a ring of ripples to refocus back into a single point, like a reversed movie of the dropped stone. But we never actually see this happening, because it is vanishingly unlikely that an outside-in ring would form spontaneously. Put it in reverse

(4096) Fort and others have previously shown that it is possible to engineer the situation so the waves turn back on themselves by kitting out the edge of your pond with sensors and wave generators, effectively recording and then playing back the waves in reverse. The process has been done with water, sound and even microwaves, and is used to produce clearer ultrasound images for medical uses.

(4097) But now the team has discovered that simply dropping the tub downwards is enough to send water waves scurrying in reverse, with no need for complex sensors. That's because the velocity of water waves depends on acceleration due to gravity. "If you do a sudden jolt, you change the effective gravity, and the velocity of the wave changes," says Fort.

(4098) The disruption essentially resets the starting point of the wave. Once normal gravity is restored, it starts to travel both backwards and forwards. "Suddenly, you freeze the wave. When you release it again, it doesn't remember that it was going one way," says Fort.

(4099) Ripples in the water returned exactly to their start point - the shape of the Eiffel Tower

(4100) Ripples in the water returned exactly to their start point – the shape of the Eiffel Tower

**Déjà vu**

(4101) To show that the backward-travelling waves converge exactly on their starting point, just as if you were rewinding time, the team conducted a series of experiments. In some, they simply created a circular ripple of waves by dipping a probe in the water, similar to dropping a stone in the pond. In others, they used wave emitters in the shape of France, the Eiffel Tower and a smiley face to create complex wave-fronts. In each case, the reversed waves converged to recreate their original shape.

(4102) The experiments aren't truly reversing time, because the waves continue to travel out again once they have converged at the central point, points out Fort. "A real time-reversal would have a sink that absorbs all the waves," he says.

(4103) Now the team hopes to conduct similar experiments with other kinds of waves, like light or sound, although it may prove more difficult to find an analogue of their time-mirror jolt in other wave mediums, he says.

(4104) "The work is spectacular," says Geoffrey Edelmann of the US Naval Research Laboratory in Washington, D.C. "They essentially change the rules that govern how the wave propagates and, in doing so, create an entirely new way to focus waves free of spatial constraints." He also thinks the technique could be extended to other kinds of waves, and perhaps have medical applications such as lithotripsy, which uses ultrasound waves to break up kidney stones. Journal reference: Nature Physics, DOI: 10.1038/nphys3810

**Time reversal and holography with spacetime transformations Vincent Bacot, Matthieu Labousse, Antonin Eddi, Mathias Fink & Emmanuel Fort Nature Physics (2016) doi: 10.1038/nphys3810**

(4105) Wave control is usually performed by (e) spatially engineering the properties of a medium.

(4106) Because time and space play similar roles (e&eb) in wave propagation, manipulating time boundaries provides (eb) a complementary approach.

(4107) Here, authors experimentally demonstrate the relevance of this concept by (e) introducing instantaneous time mirrors.

(4108) Authors show with water waves that a sudden change of (e) the effective gravity generates (eb) time-reversed waves that refocus at the source.

(4109) They generalize this concept for (e) all kinds of waves, introducing (eb) a universal framework which explains (eb) the effect of any time disruption on (e&eb) wave propagation.

(4110) Authors show that sudden changes of the medium properties generate (eb) instant wave sources that emerge (eb) instantaneously from the entire space at (eb) the time disruption.

(4111) The time-reversed waves originate from (e) these 'Cauchy sources', which are (+) the counterpart of Huygens virtual sources on (eb) a time boundary. It allows us to revisit the holographic method and introduce a new approach for wave control.

**Viewpoint: A Bird's Eye View of Circuit Photons Sebastian Schmidt, Institute for Theoretical Physics, ETH Zurich, 8093 Zurich, Switzerland June 30, 2016• Physics 9, 74 A scanning probe detects the quantum states of photons in a microwave circuit, providing the information needed for quantum simulations. Figure caption expand figure A. Houck/Princeton (deleted)**

(4112) Figure 1: Scanning defect microscopy provides a map of photons in a resonator lattice. Houck and colleagues demonstrated the technique using (e) 49 resonators (grey lines) that were coupled together to (e&eb) form a kagome lattice.

(4113) Quantum mechanics rules the dynamics of (e) light and matter.

(4114) Yet performing a quantum-mechanical simulation of a material from first principles is (=) practically impossible on a classical computer because (e) the complexity of the simulation increases (eb) exponentially with the number of particles involved.

(4115) The solution, according to Richard Feynman, was to build a machine out of (e) quantum building blocks that could directly emulate (e&eb) the material itself [1].

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| (4116) Prototypes of such quantum simulators that are based on (e) ultracold atoms, ions, photons, and superconducting microwave circuits are now available [2], with (e&eb) the latter, in particular, having attracted (e) Silicon Valley's interest. |
| (4117) The challenge with these circuit-based simulators, however, is that they are (=) 2D, which complicates the readout of their constituent elements. Andrew Houck from Princeton University, New Jersey, and colleagues have now delivered an attractive solution by developing a technique [3], called scanning defect microscopy (Fig. 1), that determines (eb) the number of photons occupying (e) each mode of a 2D microwave circuit. |
| (4118) It is this information that would serve as (=) the fundamental input and output for (e) certain quantum simulations. |
| (4119) Superconducting microwave circuits combine electronic and (e&eb) photonic degrees of freedom [4, 5]. |
| (4120) The main element of the circuit is (=) a transmission line, which is made up of (e) a central superconducting wire separated by a gap from (e) two grounded plates. |
| (4121) All of these structures are on a single plane, as if one had taken a 2D slice through (e&eb) a coaxial cable. |
| (4122) When truncated, the transmission line becomes (=) a resonator, which can host (e&eb) discrete photon modes within (eb) its gaps |
| (4123) Large lattices of resonators can be engineered in various 1D or 2D geometries by coupling (e&eb) two, three, or more resonators together via (e&eb) a capacitive interface. In many ways, photons in such devices behave similarly to electrons in a solid. |
| (4124) To make microwave circuits that can simulate (e&eb) quantum phenomena faster than a classical computer, however, resonator lattices have to be integrated with (e&eb) superconducting qubits |
| (4125) . Such qubits are controlled with (e&eb) electrical currents in a Josephson tunnel junction, and in many respects, they behave (e&eb) like artificial atoms, which couple to (e&eb) the photons in the circuit. |
| (4126) As a result, superconducting microwave circuits can be used to (e) explore the coupling between the quantum states of light and (e&eb) matter, the regime of (e) circuit quantum electrodynamics (cQED). |
| (4127) Photons in these devices often exhibit (eb) striking matter-like behavior [6, 7], providing (eb) the basis for the simulation of complex materials. |
| (4128) Such circuits can be fabricated on a substrate using (e) standard lithographic techniques, with (e&eb) qubits and resonators that are (=) hundreds of micrometers or even millimeters in size. |
| (4129) Small-scale realizations of cQED lattices have already been used for (e) proof-of-principle quantum simulations [8, 9]. |
| (4130) However, future large-scale simulations will have to be initialized and measured (eb) efficiently—steps known as (=) drive and readout. |
| (4131) This requires connecting the lattices to (e) external circuitry, such as voltage sources, circulators, and amplifiers, which typically reside outside (e) the chip. |
| (4132) So far, this connectivity has been realized through (e&eb) external transmission lines, but these elements can only couple to (e&eb) the resonators at the edges of (e) the lattice because of (e) its 2D nature. |
| (4133) Such a side view provides (eb) only limited information about (e&eb) the quantum state and spatial correlations of photons inside (eb) the bulk resonators, seriously limiting (e) the applicability of circuit QED devices for (e) quantum simulations. |

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| (4134) The probe developed by Houck and colleagues, scanning defect microscopy (SDM), solves (e) this problem by (e) moving outside the plane of the chip to provide (eb) a fuller image of photons within (eb) the chip. |
| (4135) In SDM, a small sapphire tip is positioned a few tens of (e) micrometers above the surface of one of the resonators (Fig. 1). |
| (4136) The presence of this dielectric modifies (e&eb) the evanescent electric field from the circuit photons, which extends (eb) just outside the plane of the chip. |
| (4137) Specifically, the tip increases (eb) the local capacitance of the transmission line, downshifting (e) the local resonator frequency. |
| (4138) This in turn causes (eb) a drop in the transmission of photons through (e&eb) the lattice when (e) it is driven at its nominal resonance frequency. |
| (4139) Houck's team showed that if the sapphire tip has only a weak (linear) effect on the resonator frequency, then (eb) a simple analytic relationship exists between the induced frequency shift, (e&eb) the change in transmission, and (e&eb) the local photonic occupation of each resonator. |
| (4140) This relationship provides (eb) a recipe for detecting the local photon occupancy inside (eb) the lattice. |
| (4141) First, calibrate the probe by determining (eb) the frequency shift for (e) a single isolated resonator as (=) a function of the tip's position. |
| (4142) Second, measure the transmission through (e&eb) the whole lattice with and without (e&eb) the probe. |
| (4143) Finally, use the measured transmission drop with (e&eb) the probe present to find (eb) the photonic occupation of (e) an individual resonator in (eb) the lattice. |
| (4144) A map of the photon-mode occupancies can be obtained by (e) scanning the probe across the whole lattice. |
| (4145) Houck's team used (e) SDM to image the occupancies of (e) the photon modes of (e) 49 coupled resonators arranged in (eb) a special geometry called a kagome lattice (see details in Fig. 1). |
| (4146) (To simplify the measurement, no superconducting qubits were (=) incorporated with (e&eb) the resonators.) |
| (4147) The occupancies they measured were in excellent agreement with (=) theoretical predictions, corresponding to (e&eb) a fidelity of 0.99. |
| (4148) The high fidelity is even more impressive considering the weak signal-to-noise ratio of (e) their technique (which is limited by (e) the intentionally small size of the frequency-shifted region) and (e&eb) the technical challenge of (e) positioning the probe inside (eb) a cryostat for (e) in situ measurements. |
| (4149) While SDM shares similarities with (e&eb) so-called bead-pull experiments for accelerators and scanning gate microscopy for (e) nanostructures, it is truly a (=) new tool for quantum electronics. |
| (4150) First application of SDM is likely to be (=) the simulation of 2D topological phenomena and (e&eb) geometrically frustrated lattices (of which the kagome lattice is (=) one example.) |
| (4151) Both are associated with (e&eb) macroscopically degenerate systems, which are (=) notoriously difficult to simulate on (e&eb) a classical computer. |
| (4152) Such circuits are (=) genuinely open systems (photons can be added or lost). |
| (4153) This combined with (e&eb) their inherent strong light-matter coupling, makes (eb) such quantum simulators ideal candidates for (e) the study of many-body, out-of-equilibrium phenomena, such as (=) thermalization or nonequilibrium phase transitions. |

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| (4154) SDM could also lead to (eb) the discovery of new exotic phenomena, which are (=) unknown in traditional solid-state systems. |
| (4155) In a nutshell, SDM constitutes a new and indispensible tool for (e) future experiments on quantum simulations with microwave photons. This research is published in Physical Review X. |

**Imaging Photon Lattice States by Scanning Defect Microscopy D. L. Underwood, W. E. Shanks, Andy C. Y. Li, Lamia Ateshian, Jens Koch, and A. A. Houck Phys. Rev. X 6, 021044 – Published 30 June 2016**

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| (4156) Microwave photons inside (eb) lattices of coupled resonators and superconducting qubits can exhibit (eb) surprising matter like behavior. |
| (4157) Realizing such open-system quantum simulators presents (eb) an experimental challenge and requires (e) new tools and measurement techniques. |
| (4158) Here, authors introduce scanning defect microscopy as one such tool and illustrate its use in mapping (e&eb) the normal-mode structure of microwave photons inside (eb) a 49-site kagome lattice of coplanar waveguide resonators |
| (4159) Scanning is accomplished by (e) moving a probe equipped with a sapphire tip across (e&eb) the lattice. |
| (4160) This locally perturbs (e&eb) resonator frequencies and induces (eb) shifts of the lattice resonance frequencies, which authors determine by (e) measuring the transmission spectrum. |
| (4161) From the magnitude of mode shifts, author can reconstruct (eb) photon field amplitudes at (eb) each lattice site and thus create (eb) spatial images of the photon-lattice normal modes. Received 4 December 2015 DOI:http://dx.doi.org/10.1103/PhysRevX.6.021044 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society |

| |
|---|
| |
| We assume that |
| $$D_7 = \mu_7(a_7) + (f_7)_7(P_7)_7 + (f_7)_7(Q)_7 \quad (7)$$ |
| $$n_7 = \beta_7(a_7)h_7(Q_7), \quad (3)$$ |
| Where $\mu_7$the age- dependent debit rate whit is $\mu_7(a_7)$ 0 $\forall a > 0$; $(f_7)_7$ is the density- dependent death rate with $f_7(0) = 0$, $(f_7^7)_7(P_7)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_7)_7(P_7) = (k_7)_7$ ∞; $f_7$ is the functional response of creditors to prey with $[f_7]_7(0) = 0, (f_7^7)_7(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_7)_7(Q_7) = [k_7]$ ∞; $\beta$ is the age dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^7(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0. |

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_7 + (\rho_a)_7 = -\mu_7(a_7) + (f_7)_7 P_7) + \big(f_7)_7(Q_7)\big)(\rho_7)_7 \quad (4)$$

$$\rho_7(t_7, 0) = h_7(Q_7) \int_0^\infty \beta_7 (a_7)\rho_7(t, a_7)da, \quad (5)$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + g_7(P_7)Q_7 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_7$ and $Q_7$ of (4-6) satisfy

$$0 \leq \rho_7(t, a_7) \leq \rho_7^*(t, a_7),$$

$$0 \leq Q_7(t) \leq Q_7^*(t),$$

Where $\rho_7^*(t, a_7)$ and $Q_7^*(t)$, are solutions of

$$(\rho_t^*)_7 + (\rho_a^*)_7 = -(\mu_7(a_7) + (f_7)_7(P_7^*) + (f_7)_7 Q_7^*))\, \rho_7^*, \quad (7)$$

$$\rho_7^*(t, 0) = \int_0^\infty \beta_7(a_7)\, \rho_7^*(t, a_7)da_7 \quad (8)$$

$$\frac{dQ_7^*}{dt} = -b_7\, Q_7^* + g_7(P_7^*)Q_7^*, \quad (9)$$

With $P_7^*(t) = \int_0^\infty \rho_7^*(t, a_7)da_7$

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From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_7^*(t)}{dt} = P_7^*(t)\big(r_7 - (f_7)(P_7^*) - (f_7)_7(Q^*)\big), \quad (10)$$

$$\frac{dQ_7^*(t)}{dt} = Q_7^*(t)(-b_7 + g_7(P_7^*)), \quad (11)$$

Where $r$ is the real solution of the equation

$$\int_0^\infty \beta_7(a)e^{-\int_0^a \mu_7(\tau)dt}e^{-ra}\,da_7 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_7^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(t)dt}\,da_7 > 1, \quad (H1)$$

Which guarantees $r_7 > 0$.

*Remark.* If $\beta_7(a) = [\beta_0]_7 a_7 e^{-\alpha_7 a_7}, \mu_7(a_7) \equiv \mu_7$ then the solution of (12) is

$$r_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_7]_7(P_7) = 0, [f_7](Q_7) = \lambda_7 Q_7\ h_7(Q_7) = 1, \quad and\ \ g_7(P_7) = c_7 P_7$$

$$[\rho_7]_7 + [\rho_a]_7 = -(\mu_7 + \lambda_7 Q_7)\rho_7$$

$$\rho_7(t,0) = \int_0^\infty [\beta_0]_7 a_7 e^{-\alpha_7 a_7} \rho_7(t,a_7)\, da_7$$

$$\frac{dQ_7}{dt} = -b_7 Q_7 + c_7 P_7 Q_7$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_7}{dt} = \omega_7 P_7 - \lambda_7 P_7 Q_7$$

$$\frac{dQ_7}{dt} = b_7 Q_7 - c_7 P_7 Q_7$$

Where $\omega_7 = \sqrt{\beta_7} - \alpha_7 - \mu_7$

## 69. THE EQUILIBRIUM $(\rho_0)_7(a_7), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_7(a_7), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_7 + (\rho_a)_7 = -\big(\mu_7(a_7) + (f_7)_7(P_7)\big)\rho,$$

$$\rho_7(t,0)$$

$$= \int_0^\infty \beta_7\,(a_7)\rho_7(t,a_7)da_7 \qquad (13)$$

The solution of (13) has the form of

$$\rho_7(t,a_7) = Ce_7^{r(t-a_7)} - \int_0^a (\tau)d\tau - \int_0^7 (f_7)_7\big(P_7(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_7(t) = C_7 e^{r_7 t - \int_0^t (f_7)_7 (P_7(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_7^{-7}(r)$ if $k_7 > r$. This is equivalent to

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{(k_7)_7 a_7} da_7 < 1.$$

In fact, if we define

$$R_7(s_7) = \int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau)d\tau} e^{-s_7 a_7}\, da_7$$

$R(s)$ is a decreasing function. $R_7(r_7) = 1$ And $R_7(k_7)_7 < 1$ imply $r_7 < (k_7)_7$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_7(a_7), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{k_7 a_7} da_7$$

$$< 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_7 \equiv \int_0^\infty (\rho_0)_{77}(a_7)da_7$ *satisfies*

$$\int_0^{\infty} \beta_7(a_7) e^{-\int_0^a \mu_7(\tau)d\tau} e^{-(f_7)_7(P_0)a_7} da$$

$$= 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_7(a_7), 0)$

Let $u_7(t, a_7) = \rho(t, a_7) - \rho_0(a_7)$ Then

$$N_7(t) \equiv \int_0^{\infty} u_7(t, a_7) da_7 = \int_0^{\infty} \rho_7(t, a_7) da_7 - \int_0^{\infty} (\rho_0)_7(a_7) da = P_7(t) - P_0)_7$$

And

$$(u_7)_7 + (u_a)_7 = (\rho_t)_7 + (\rho_a)_7 - (\rho_0^7)_7(a_7)$$

$$= -\mu_7(a_7) + (f_7)_7(P_7) + (f_7)_7(Q_7))(u_7 + (\rho_0)_7) + \mu_7(a_7)(\rho_0)_7 + (f_7)_7(P_0)_7(\rho_0)_7$$

$$= -\mu_7(a_7)u_7 - (f_7)_7(P_7) + (f_7)_7(Q_7))\, u_7 - (\rho_0)_7\, (f_7)_7(P_7) - (f_7)_7(P_0)_7$$

$$+ (f_7)_7(Q_7)), \qquad (15)$$

Where $(\rho_0^7)_7(a_7)$ denotes $d(\rho_0)_7(a_7)/da_7$

Expansion $(f_7)_7(P_7)$ and $(f_7)_7(Q_7)$ at $(P_0)_7$ and 0 yield respectively

$$(f_7)_7(P_7) = (f_7)_7((P_0)_7 + N_7)$$

$$= (f_7)_7(P_0)_7 + (f_7^7)_7(P_0)_7 N_7 + o_7(N_7), \qquad (16)$$

$$(f_7)_7(Q_7)$$

$$= (f_7^7)_7(0)Q + o_7(Q_7) \qquad (17)$$

Write $f_7(P_0) = f_{70}, f_7'(P_0) = f_{70}'$, and $f_7'(0) = f_{70}'$ and substitute (16) and (17) into (15) to obtain

$$(u_7)_7 + (u_a)_7 \approx -(\mu_7(a_7) + (f_{70})_7)u - (f_{70}^7)_7(\rho_0)_7(N_7) - (f_{70}^7)_7(\rho_0)_7(Q_7)$$

$$u_7(t, 0) = \rho_7(t, 0)(\rho_0)_7(0)$$

$$= h_7(Q_7)\int_0^{\infty} \beta_7(a_7)(u_7(t,a) + (\rho_0)_7(a_7)da_7 - (\rho_0)_7\,(0)$$

$$= h_7(Q_7)\int_0^{\infty} \beta_7(a_7)(u_7(t,a_7)da_7 + (h_7(Q_7) - 1)(\rho_0)_7\,(0) \qquad (18)$$

Expanding

$$h_7(Q_7) = 1 + \ (h^7)_7(0)Q_7 + 0(Q_7) \equiv 1 + h_0^{'}Q_7 + 0Q_7$$

And substituting it into (18), we arrive at

$$u_7(t,0) \approx \int_0^{\infty} \beta_7\,(a_7)u_7(t,a_7)da_7 + (h_{a_7}^{'})(\rho_0)_7Q_7$$





The linearization is the following :

$$[u_7] + [u_7] = -(\mu_7(a_7) + [t_{70}]_7u_7 - [\alpha_7]_7(a_7)N_7 - [\alpha_7]_7(a_7)Q_7$$

$$u_7(t,0) = \int_0^{\infty} \beta_7(a_7)u_7(t,a_7)da_7 + [\xi_7]Q_7 \qquad (19)$$

$$\frac{dQ_7}{dt} = -b_7Q_7$$

where

$$[\alpha_7]_7(a_7) = [t_7^7a]_7[p,o]_7(a_7)d2(a_7) = [f_{70}^7]_7(\rho\ o\ )_7(a_7)\xi = [\mathrm{h\ o}^7]_7[\rho\mathrm{o}]_7(\mathrm{o}) \text{ and } \mathrm{b}_7 = \mathrm{b}_7 - g_7(P_7)$$

If $g_7(P_0) > b$, clearly, $(\rho_0)_7(a_7), 0)$ is unstable and the predator population is persistent. Assume $g_7(P_0)_7 < b$.

Then, $Q_7(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_7 + u_a = -(\mu_7(a_7) + (f_{70})_7)u_7 - (\alpha_7)_7(a_7)N_7$$

$$= -(\mu_7(a_7) + f_{70})u_7 - \alpha_7(a_7)\int_0^\infty u_7(t, a_7)da_7$$

$$u_7(t, 0) = \int_0^\infty \beta_7\,(a_7)u_7(t, a_7)da_7$$

Using the method developed by Cushing [2, 3], we suppose

$$u_7(t, a_7) = \tilde{u}_7(a_7)^{e^{\gamma(t-a_7)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_7}\,da_7$. Then,

A solution of the first differential equation for $\tilde{u}(a_7)$ gives,

$$\tilde{u}_7(a_7) = e^{-\int_0^a \mu_7(\tau_7)d\tau - (f_{70})_7 a}\left(\tilde{u}_7(0) - w\int_0^a e^{\int_0^\tau \mu_7(s_7)ds + (f_{70})}\tau_7\,(\alpha_7)_7(\tau_7)e^{\gamma_7\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_7(a_7)e^{-\gamma_7 a_7 - (f_{70})_7 a_7 - \int_0^a \mu_7(\tau_7)d\tau}\right)\tilde{u}_7(0)$$

$$+w\int_0^\infty \beta_7(a_7)e^{-\gamma_7 a_7 - (f_{70})_7 a_7 - \int_0^a \mu_7(\tau_7)d\tau}\int_0^a (\alpha_7)_7(\tau_7)e^{-\gamma_7\tau_7}e^{(f_{70})_7\tau_7 + \int_0^\tau \mu_7(s_7)ds}d\tau_7\,da_7$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of $w$ yields

$$\left(\int_0^{\infty} e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau} da\right)(\tilde{u}_7(0))_7$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(\tau_7)d\tau}\int_0^a e^{(f_{70})_7\tau_7+\int_0^\tau \mu_7(s_7)ds+y\tau}\alpha_7(\tau)d\tau\, da\right)w=0. \quad (23)$$

Define $Y(\gamma,a)=e^{-\gamma_7 a_7-(f_{70})_7 a_7-\int_0^a \mu_7(t_7)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)da_7\right)u_7(0)$$

$$+w\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\, da_7=0, \quad (24)$$

$$\left(\int_0^{\infty}Y_7(\gamma_7,a_7)da_7\right)\tilde{u}_7(0)$$

$$+\left(-1-\int_0^{\infty}Y_7(\gamma_7,a_7)\int_0^a(\alpha_7)_7(\tau)_7(Y^{-7})_7(\gamma_7,\tau_7)d\tau_7\, da_7\right)w=0.$$

Equation (20) has a solution of the form of $u_7(t,a_7)=u_7(a_7)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_7(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_7(\gamma_7,a_7)=\int_0^a(\rho_0)_7(\tau_7)Y^{-7}(\gamma_7,a_7)d\tau_7$. We obtain the characteristic equation

$$f_{70}^{l}{}_{7}\left(\int_0^{\infty} Y_7(\gamma_7,a_7)da_7\right)\left(\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$+\left(1-\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)da\right)\left(1+(f_{70}^{l})_7\int_0^{\infty}Y_7(\gamma_7,a_7)I_7(\gamma_7,a_7)da_7\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_7(a_7)=(\rho_0)_7(0)e^{-\int_0^a \mu_7(\tau_7)d\tau_7-(f_{70})_7a_7}$$

$$I_7(\gamma_7,a_7)=\int_0^a(\rho_0)_7(\tau_7)e^{\gamma_7\tau_7+(f_{70})_7\tau+\int_0^{\tau}\mu_7(s_7)ds_7}d\tau$$

$$=\frac{(\rho_0)_7(0)}{\gamma_7}(e^{\gamma_7a_7}-1)$$

Let

$$E_7(a_7)=e^{-(f_{70})_7a_7-\int_0^a\mu_7(\tau_7)d\tau_7}$$

Then $[Y_7]_7[\gamma_7]_7,a_7)=E_7(a_7)e^{-\gamma_7a_7}$. by a simple calculation,

$$\int_0^{\infty}\beta_7(a_7)Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da=\frac{[\rho_0]_7(0)}{\gamma}\left(1-\int_0^{\infty}\beta_7(a_7)E_7(a_7)e^{-\gamma_7a_7}\ da\right);$$

$$\int_0^{\infty}Y_7(\gamma_7,a_7)I(\gamma_7,a_7)da=\frac{[\rho_0]_7(0)}{\gamma_7}\left(\int_0^{\infty}E_7(a_7)e^{-\gamma_7a_7}\ da_7\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_7(0)}{\gamma_7}[f_{70}^7]\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\,(\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7)$$

$$+\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right).$$

$$\left(1+\frac{(\rho_0)_7(0)}{\gamma_7}(f_{70}^{'})_7\left(\int_0^\infty E_7(a_7)da_7-\int_0^\infty E_7(a_7)e^{-\gamma_7 a_7}da_7\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7\right)(1+\frac{(f_{70}^{'})_7}{\gamma_7}(\rho_0)_7(0)\int_0^\infty E_7(a_7)da_7)=0.$$

We need $\tilde{u}(a) \not\equiv 0$, namely $w \neq 0$. This necessarily leads, from (24), to

$$1-\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)da_7 = 1-\int_0^\infty \beta_7(a_7)E_7(a_7)e^{-\gamma_7 a_7}da_7 \neq 0,$$

Since $\int_0^\infty \beta_7(a_7)Y_7(\gamma_7 a_7)\int_0^a \alpha_7(\tau_7)Y^{-7}(\gamma_7,\tau_7)d\tau_7 da_7 > 0$. Therefore, the characteristic solution is

$$\gamma_7 = -(f_{70}^{'})_7(\rho_0)_7\int_0^\infty E_7(a_7)da_7.$$

Since $f_7'(P_0)_7 > 0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_7), 0)$ is asymptotically stable if $g(P_0) < b$. It is unstable if $g(P_0) > b$.

## 3. COEXISTENCE OF DEBIT AND CREDIT TYPES

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_7), 0)$ which is related to the extinction of predators and to The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_7 + (\mu_7(a_7) + [f_7]_7(P_7^*) + [f_7]_7(Q_7^*))P_7^* = 0, \quad (76)$$

$$p_7^*(0) = h_7(Q_7^*)\int_0^\infty \beta_7(a_7)p_7^*(a_J)p_7^*(a_7)da_7 \quad (77)$$

$$g_7(P_7^*) - b_7 = 0, \quad (78)$$

Where $P_7^* = \int_0^\infty p_7^*(a_7)\, da_7$

Equation (28) yields $P^* = g^{-7}(b)$,

From (26),

$$p_7^*(a) = p_7^*(0)e^{\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(Q_7^*)a_7}$$

Substituting into (27), we need to solve

$$\left(1 - h_7(Q_7^*)\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(Q_7^*)a_7}da_7\right)p_7^*(0) = 0$$

For $Q_7^*$

Fix $p_7^*(0)$ and $P_7^*$ Let

$$H_7(z_7) = 1 - h_7(z_7)\int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(z_7)a_7}da_{,7}$$

And assume that

$$[k_3]_7\int_0^\infty \beta_7(a_7)e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(g_7^{-7}(b_7))a_7 - [k_7]_7 a_7}\, da_7 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_7(z_7)] > 0.$$

Clearly,

$$\frac{dH_7}{dz_7} = h_7(z)(\,[f_7']_7(z_7)\,)\int_0^\infty a_7\beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - f_7(P_7^*)a_7 - f_7(z_7)a_7}da_{7}$$

$-h'(z_7)\int_0^\infty \beta_7(a_7)\, e^{\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7 - [f_7]_7(z_7)a_7}\quad da_7$

(See author's It is a debit-credit world)

By the assumptions on $\left[f_7^7\right]_7 and\ [h^7]^7 dH_7/dz > 0$. Hence, if

$$H_7(0) = 1 - \int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7(P_7^*)a_7} da_7 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_7(0) < 0$ follows if $g_7^-(b_7) < [f_7^{-7}]_7(r_7)$, where $r_7$ is the solution of (12).

Indeed, since

$$R_7(s_7) = \int_0^\infty \beta_7(a_7)\, e^{-\int_0^a \mu_7(\tau_7)d\tau_7 e^{-y\frac{a}{7}}} da_7$$

Is a decreasing function of $s_7$ when $P_7^* = [g^{-7}]_7(b_7) < [f_7^{-7}]_7(r_7)$, that is $[f_7]_7(P_7^*) < r$, $R\left(f_7(P^*)\right) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

**THEOREM 3**

Assume that the hypothesis (H3) holds and $R(f_7(g^{-7}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_7(a_7) e^{-\int_0^a \mu_7(\tau_7)d\tau_7 - [f_7]_7\left(g_7^{-7}(b)\right)a_7} > 1. \qquad (H4)$$

Then the system (4)-(6) has a unique positive equilibrium $(P_7^*(a_7), Q_7^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_7^*(a_7), Q_7^*)$ is more complicated. We again consider perturbations $m_7(t_7 a_7) = p_7(t_7, a_7) - p_7^*(a_7)$ and $W_7(t) = Q_7(t) - Q^*$.Then

$M_7(t) \equiv \int_0^\infty m_7(t,a_7)da_7 = P_7(t) - P_7^*,$

$$[m_7]_7 + [m_a]_7 = [p_7]_7 + [p_a]_7 - [p^*]_7(a_7) \approx -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7$$

$$-[f_{7*}^7]_7 p_7^*(a_7)M - [f_{7*}^7]_7 p_7^*(\mathrm{a})W_7,$$

$$m_7(t,0)p_7(t,0) - p_7^*(0) \approx h_* \int_0^\infty \beta_7(a_7)m_7(t,a_7)\,da + \frac{[h_*^7]_7}{[h_*]_7}\, p_7^*(0)W_7$$

$$\frac{dW_7}{dt} = \frac{dQ_7}{dt} \approx [g_*^7]_7 Q_7^* M,_7$$

Where $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}]_7$ denoting $[f_7(P^*)]_7$, $[f_{7*}']_7$ denoting $[f'_7(P^*)]_7[f_{7*}']_7$ denoting

$[f_7]_7$ $(Q_7^*,$ $h_{7}^*$ denoting $h_*'$ denoting $h_7'$ $(Q_7^*)$, and $[g_*']_7$

Denoting $[g']_7(P_7^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$[m_7]_7 + [m_a]_7 = -(\mu_7(a_7) + [f_{7*}]_7 + [f_{7*}]_7)m_7 - [f'_{7*}]_7[p_7^*](a_7)M_7 - [f_{7*}']_7 p_7^*(a_7)W_7$

$m_7(t,0) = [h_*]_7 \int_0^\infty \beta_7(a_7)m_7(t,a_7)da_7 + \frac{[h_*^7]_7}{[h^*]_7}[p^*]_7(0)\,W,_7$ (30)

$$\frac{dW_7}{dt} = [g_*^7]_7[Q_7^*]M_7$$

Integrating $[m_7]_7(t,a_7)$ with respect to a from 0 to ∞, we have

$$\frac{dM_7(t)}{dt} = \int_0^\infty [mt]_7(t, a_7)da_7$$

$$= -\int_0^\infty [m_a]_7\,(t, a_7)da_7 - \int_0^\infty \mu(a_7)m(t, a_7)da_7$$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f^{'}{}_{7*}]_7 \int_0^\infty p_7^*(a_7)da_7\right)M_7 - [f_{7*}^7]_7 \int_0^\infty p_7^*(a_7)da_7 W_7$$

$$= m_7(t, 0) - \int_0^\infty \mu_7(a_7)\, m_7(t, a_7)da_7$$

$$-\left([f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 - [f_{7*}^{'}]_7 P_7^* W_{.7}$$

From (30), we arrive at the equations

$$\frac{dW_7(t)}{dt} = \int_0^\infty ([h_*]_7\beta_7(a_7) - \mu_7(a_7))m_7(t, a_7)da_7$$

$$-\left([f_{7*}] + [f_{7*}]_7 + [f_{7*}^{'}]_7 P_7^*\right)M_7 + \left(\frac{[h_*^{'}]_7}{[h_*]^7}\, P_7^*(0) - f_{7*}^{'} P_7^*\right)W_7,$$

(31)

$$\frac{dW_7(t)}{dt}[g^{'}{}_*]_7\, Q_7^*\, M_{.7}$$

By the mean value theorem,

$$\int_0^\infty ([h_{"*}]_7\beta_7(a_7) - \mu_7(a_7))m_7(t, a_7)da_7 = ([h_*]_7\beta_7(\theta_7) - \mu_7(\theta_7))M,_7$$

Where $0< \theta_7 = \theta_7(t) < \infty$.

And if a $_7^7$ [a $_7^7$]2<4 [a $_7$]2<4 a $_7$e,

Re $\lambda_7 [\lambda_7]_7 = \text{Re}\lambda_7 [\lambda_7]_7 = -\frac{7}{7} a_7$[a2 ] $_7 < 0$.

Let ⊿=max$_{7<u<\infty}$ |h $_*$ [h $_*$] $_7\beta_7(a_7) - \mu_7(a_7)$|. Then|$n_7$(t)| 2⊿∀t 0. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium($\rho^*$(a), $\boldsymbol{Q}^*$) is asymptotically stable if either a $_7^7$ <4 a $_7$ e and [a $_7$] > 1⊿,or [a $_7^7$] 24[a $_7$]e and (a $_7 - \sqrt{(\text{a}_7^7)_7 - 4(\text{a}_7)_7 \text{e})} > 1⊿$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{\text{a}}{\text{a}+7}\right)$ and $\beta_7(a_7) =$ ae-$^{\text{ma}}$ Then ⊿=[$\mu_0$] + 2 Simple calculation leads to

$$\int_0^\infty \beta_7(a_7)\text{e}^{-\int_0^{\text{a}} \mu_7(\tau_7) + d\tau_7}\text{da} = \frac{1}{(\text{m} + \mu_0 + 1)^7} + \frac{1}{(\text{m} + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_7 + [\mu_0]$ 0.5. Assume$[f_7](x_7)= K_7x_7$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_7(x_7)= [0x_7]/(2 + x_7)$; the functional responses $[[f_7](x) = [0]\, x_7/(1 + x_7)$and $h_7(x_7)$=(1+$\omega_7x_7$)/(1 + $x_7$)(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_7 + [\mu_0] + [\text{K}][P_0])}{[m_7]_7 + (\mu_0)_7 + [\text{K}][P_0])_7^3\ [\text{K}]_7[P_0]_7^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1+P_0) < b_7$ i e $[P_0] < b/(\delta_7 - b_7)$then the equilibrium$(P_0)$ $[P_0]_7(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_7$

$[f_{7*}] = K_7\ \mathrm{P}^*\ [f_{7*}] = \frac{\sigma \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*} \quad [h_*] = \frac{7+\omega_7 \mathrm{Q}_7^*}{7+\mathrm{Q}_7^*}$

$[f_{7*}] = K_7\ [\mathrm{f}_{7*}^{\ 7}] = \frac{\sigma_7}{(7+\mathrm{Q}_7^*)^7} [\ \mathrm{g}_*^{\ 7}] = \frac{\sigma_7}{(7+\mathrm{P}_7^*)^7} \quad [\mathrm{h}_*^{\ 7}] = \frac{\omega_7 - 7}{(7+\mathrm{Q}_7^*)^7}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta,$ and $\omega$ need to satisfy

$$\frac{\omega_7}{(m + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma)^7} + \frac{1\omega_7}{([m_7]_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 + b_7} + \sigma_7)^3} < 1$$

And

$$\frac{1}{(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^7} + \frac{1\omega_7}{(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $\mathrm{Q}^*$ which can be solved from the following equation:

$(m_7 + [\mu_0]_7 + 1 + \frac{K_7 b_7}{\delta_7 - b_7} + (m_7 + [\mu_0]_7 + \frac{K_7 b_7}{\delta_7 - b_7} + \sigma_7)\ \mathrm{Q}_7^*)^3$

$= m_7 + [\mu_0]_7 + 3 + \frac{K_7 b_7}{\delta_7 - b_7}(1 + \mathrm{Q}^*)(1 + \omega_7 \mathrm{Q}_7^*$

Now

$$\left[\mathrm{a}\begin{matrix}1\\1\end{matrix}\right]-4\left[\mathrm{a}_{7}\right]\mathrm{e}=\left(\frac{1K_7 b_7}{\delta_7-b_7}\right)^{7}+\frac{4K_7 b_7\sigma_7\,\mathrm{Q}_{7}^{*}}{(\delta_7-b_7)(7+\mathrm{Q}_{7}^{*})}+\frac{\sigma_{7}^{7}\mathrm{Q}_{7}^{*7}}{\left(1+\mathrm{Q}_{7}^{*}\right)^{7}}$$

$$\frac{4\,b_7\sigma_7\,(\delta_{7}-\mathrm{b})\mathrm{Q}^{*}}{\delta_{7}\,(1+\mathrm{Q}_{7}^{*})^{7}}-\frac{4\,(\delta_{7}-b_7)^{7}\,(1-\omega_7)\mathrm{Q}_{7}^{*}}{\delta_{7}\left(1+\omega_7\mathrm{Q}_{7}^{*}\right)(1+\mathrm{Q}_{7}^{*})}\mathrm{P}^{*}(0)$$

And

$$\Delta=\lim_{0<a<\infty}|[h_*]_{7}\,ae_{7}^{mu}\;\;[\mu_0]_{7}-\frac{a_7}{a_7+7}$$

Hence if $[\mathrm{a}_{7}^{7}>4(a_7)_{7}\mathrm{e}$ and

In the following are given some more modular structures that are applied with the model and the concomitant results derived thereof:

Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments

(1) We have an intuitive understanding of how the world works. Our intuition works well in our everyday lives, but it doesn't quite match (e&eb) reality.

(2) One of the basic assumptions authors make about objects is that they are (=) objectively real.

(3) For example, when you leave your apartment in the morning, you imagine your house is right where you left it as you go about your day. Barring some calamity, you expect to find your apartment just as you left it when (e) you return in the evening.

(4) This idea is known as Macrorealism, in that a macroscopic object such as your house is still there even when (e) you're not looking at it

(5) In scientific terms we would say the state of an object exists independent of (e) an observer. While this seems obviously true, we can test the idea experimentally.

(6) To show this, let's consider an experiment. Suppose we have a box that contains a coin. At any time we can open the box and (eb) see whether it's heads or tails.

(7) The catch is (=) that the box has to follow a certain rule. Whatever you observe initially, (eb) you must observe the same thing ten seconds later.

(8) If you open the box and see heads, then (eb) ten seconds later when you open the box you must also see heads.

$$[a_7] = \frac{1K_7 b_7}{\delta_7 - b_7} + \frac{\sigma_7 \, \mathrm{Q}_7^*}{1 + \mathrm{Q}_7^*} > 1\vartriangle$$

This equation is important in its consequence in respect of following cases for which the model is applied:

**Why Neutrinos Aren't "Real" In Quantum Mechanics by Brian Koberlein29 August 20160 Comments**

(43) For the experiment we'll open the box, then open it five seconds later, then finally open it again at ten seconds. We know the first and last observations have to agree, but what about the observation in the middle? It depends upon what the box does with (e&eb) the coin.

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| (44) One possibility is that the box does nothing, so if you see tails the first time (eb), you will see tails both other times. |
| (45) Another possibility is that each time you open the box (eb) it flips the coin when you close it. |
| (46) So the first time you see heads, (eb) the second time tails, and (eb) the third time heads again. |
| (47) But they could also make a tricky box that flips (e&eb) the coin at seemingly random times, but does so in such a way that every ten seconds it repeats (eb) the pattern. In that case, authors would always see the same thing the first and last time, but the middle observation would seem (eb) random. |
| (48) If Macrorealism is true, then (eb) middle observations being always the same as (=) the first and last observation, always opposite or always random are (=) the only possibilities for the experiment |

**Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry Download as PDF (Size: 62KB) HTML PP. 50-53 DOI: 10.4236/jqis.2011.12007 5,508 Downloads 14,712 Views Citations Author(s) Leave a comment Mohamed S. El Naschie**

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| (43) Building upon the pioneering work of J. Bell [1] and an incredible result due to L. Hardy [2] it was shown (eb) that the probability of quantum entanglement of two particles is (=) a maximum of 9.0169945 percent [2]. |
| (44) This happens to be exactly the golden mean to (e) the power of five (?5) [3-7]. |
| (45) Although it has gone largely unnoticed for a long time, this result was essentially established independently in a much wider context by (e) the present author almost two decades ago [3-6]. |
| (46) The present work gives two fundamentally different derivations of (e) Hardy's beautiful result leading to (eb) precisely the same general conclusion, namely that by virtue of (e) the zero measure of the underlying |
| (47) Cantorian-fractal spacetime geometry the notion of (e) spatial separability in quantum physics is (=) devoid of any meaning [7]. The first derivation is purely logical and uses (e) a probability theory which combines the discrete with (e&eb) the continuum. |
| (48) The second derivation is purely geometrical and topological using (e) the fundamental equations of a theory developed by the author and his collaborators frequently referred to as E-infinity or Cantorian spacetime theory [3-7]. KEYWORDS Hardy's Quantum Entanglement, Golden mean, Cantor sets, Fractal Spacetime, E-Infinity Theory, Quantum Mechanics, J. S. Bell Cite this paper M. El Naschie, "Quantum Entanglement as a Consequence of a Cantorian Micro Spacetime Geometry," Journal of Quantum Information Science, Vol. 1 No. 2, 2011, pp. 50-53. Doi: 10.4236/jqis.2011.12007. |

**Chaos, Solitons & Fractals Volume 28, Issue 4, May 2006, Pages 857–878 Sir Hermann Bondi 1919–2005 Cantorian spacetime and Hilbert space: Part I—Foundations G. Iovane http://dx.doi.org/10.1016/j.chaos.2005.08.074**

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| (71) Authors are going to show the link between the ε(∞) Cantorian space and( e&eb) the Hilbert spaces H(∞). |
| (72) In particular, El Naschie's ε (∞) is (=) a physical spacetime, i.e. an infinite dimensional fractal space, where (e) time is spacialized and the transfinite nature manifests (eb) itself. |
| (73) El Naschie's Cantorian spacetime is an arena where (e) the physics laws appear at each scale in a self-similar way linked to (e&eb) the resolution of the act of observation. |
| (74) By contrast the Hilbert space H (∞) is (=) a mathematical support, which describes (eb) the interaction between the observer and (e&eb) the dynamical system under (e&eb) measurement. |

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| (75) The present formulation, which is based on (e) the non-classical Cantorian geometry and topology of spacetime, automatically solves (e) the paradoxical outcome of (e) the two-slit experiment and the so-called particle–wave duality. |
| (76) In particular, measurement (i.e. the observation) is equivalent to (=) a projection of ε(∞) in the Hilbert space built on 3 + 1 Euclidean spacetime. |
| (77) Consequently, the wave–particle duality becomes (=) a mere natural consequence of (e) conducting an experiment in (eb) a spacetime with (e&eb) non-classical topological and (e&eb) geometrical structures, while (e) observing and taking measurements in (eb) a classical smooth 3 + 1 Euclidean spacetime. |
| (78) In other words, the experimental fact that a wave–particle duality exists is (=) an indirect confirmation of the existence of ε(∞) and a property of (e) the quantum–classical interface. |
| (79) Another direct consequence of the fact that real spacetime is) =) infinite dimensional hierarchical ε(∞) is (=) the existence of scaling law R(N), introduced by the author, which generalizes the Compton wavelength. |
| (80) It gives an answer to the problem of (e) segregation of matter at different scales, and shows (eb) the role of fundamental constants such as (=) the speed of light and Plank's constant h in the fundamental lengths scale without (e) invoking the principles of quantum mechanics. Copyright © 2005 Elsevier Ltd. All rights reserved |

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**Chaos, Solitons & Fractals Volume 14, Issue 9, December 2002, Pages 1341–1351 Fractal strings as an alternative justification for El Naschie's Cantorian spacetime and the fine structure constant Carlos Castro, http://dx.doi.org/10.1016/S0960-0779(02)00110-8**

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| (57) Beginning with the most general fractal strings/sprays construction recently expounded in the book by Lapidus and Frankenhuysen, it is shown how the complexified extension of El Naschie's cantorian-fractal spacetime model belongs to (e&eb) a very special class of families of fractal strings/sprays whose scaling ratios are given by (e) suitable pinary (pinary, p prime) powers of the Golden Mean. |
| (58) They then proceed to show why the logarithmic periodicity laws in (eb) Nature are direct physical consequences of (e) the complex dimensions associated with (e&eb) these fractal strings/sprays. |
| (59) Authors proceed with a discussion on quasi-crystals with (e&eb) p-adic internal symmetries, von Neumann's Continuous Geometry, the role of wild topology in (e&eb) fractal strings/sprays, (e&eb) the Banach-Tarski paradox, and tesselations of (e) the hyperbolic plane, quark confinement and (e&eb) the Mersenne-prime hierarchy of (e) bit-string physics in determining the fundamental physical constants in Nature. Fax: +1-914-965-3397 Copyright © 2002 Published by Elsevier Ltd. About ScienceDirectRemote access Shopping cart Contact and support Terms and conditions Privacy policy Cookies are used by this site. For more information, visit the cookies page. Copyright © 2016 Elsevier B.V. or its licensors or contributors. Science Direct ® is a registered trademark of Elsevier B.V. |

**Chaos, Solitons & Fractals Volume 21, Issue 2, July 2004, Pages 261–271 On Cantorian spacetime over number systems with division by zero Jakub Czajko, http://dx.doi.org/10.1016/j.chaos.2003.12.046**

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| (60) Division by zero is proposed in order to open (eb) new horizons for (e) abstract mathematics and physics. |
| (61) Although the formerly forbidden operation may be of (e) little practical value, the feasibility of unrestricted operations is needed for (e) meaningful deployment of linear vector spaces, whose duals are (=) inverse transposes of their primary vector spaces. |
| (62) Both quantum mechanics and (e&eb) pan-geometry of multispatial hyperspace operate on (e&eb) vectors that belong to (e&eb) mutually dual linear vector spaces. |
| (63) The division by zero also may remove (e) many years old paradox related to (e&eb) multiplication of zeros and makes it possible to introduce (eb) the concept of bigroup as (=) an additive and multiplicative group at the same time. |
| (64) This operational extension to (e&eb) algebraic groups makes thus (eb) an infinite-dimensional, fractal Cantorian spacetime (=) the preferred, invertible, operationally unrestricted, abstract mathematical infrastructure for physics. Tel.: +1-925-681-2058 Copyright © 2003 Elsevier Science Ltd. All rights reserved. |

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The nontrivial equilibrium is stable.

Some of the examples in which stability analysis is of seminal importance are: These are the system to which the

model becomes applicable

**Observation of a Dissipation-Induced Classical to Quantum Transition J. Raftery, D. Sadri, S. Schmidt, H. E. Türeci, and A. A. Houck Phys. Rev. X 4, 031043 – Published 8 September 2014**

(699)Here, authors report the experimental observation of (e) a dynamical quantum phase transition in (eb) a strongly interacting open photonic system.

(700) The system studied, comprising a Jaynes-Cummings dimer realized on (eb) a superconducting circuit platform, exhibits (eb) a dissipation-driven localization transition.

(701)Signatures of the transition in the homodyne signal and photon number reveal (eb) this transition to be from (e) a regime of classical oscillations into (e&eb) a macroscopically self-trapped state manifesting revivals, (=) a fundamentally quantum phenomenon.

(702) This experiment also demonstrates (eb) a small-scale realization of a new class of quantum simulator, whose well-controlled (e&eb) coherent and dissipative dynamics is suited to (e) the study of quantum many-body phenomena out of (e) equilibrium. Received 10 December 2013 DOI:http://dx.doi.org/10.1103/PhysRevX.4.031043 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**Viewpoint: Fermionic Vortices Find their Dual John McGreevy, Department of Physics, University of California at San Diego, La Jolla, CA 92093-0319, USA June 27, 2016• Physics 9, 72Theoretical work reveals a surprising relationship between the physics of fermionic vortices and quantum electrodynamics. Figure caption expand figure APS/Alan Stonebraker Figure 1: The action of discrete symmetries on the fermionic particle (top) and its dual fermionic vortex (bottom).**

(703)Dirac fermions in 2+1 dimensions (two spatial dimensions plus time) have been found in (eb) graphene and on (e&eb) the surfaces of topological insulators (TIs).

(704) A trio of papers now shows that these particles can also be found in (eb) a dual description of quantum electrodynamics (QED) [1–3].

(705)This deceptively simple statement joins (e&eb) together several disparate lines of progress in theoretical physics: interacting topological insulators, (e&eb) anomalies in quantum field theory, (e&eb) spin liquids, and (e&eb) compressible Quantum Hall States.

(706)Humans are not very good at understanding interacting many-body systems.

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| (707)We understand the behavior of free particles and can sometimes use (e) perturbation theory to (e) incorporate weak interactions between (e&eb) them. |
| (708)But what happens when we couple (e&eb) the particles strongly? |
| (709)To reveal the interesting collective phenomena that can result, we've had to rely largely on (e) experiments. |
| (710)Examples include (e) superconductivity (e&eb) superfluidity (e&eb) the integer and (e&eb) fractional quantum Hall effects (e&eb) the fractionalization of charge and spin, and (e&eb) the wetness of water. |
| (711)The fractional quantum Hall effect (FQHE) is a physical phenomenon in which the Hall conductance of 2D electrons shows (eb) precisely quantised plateaus at (eb) fractional values of (e) {\displaystyle e^{2}/h} e^{2}/h. |
| (712) It is a property of (e) a collective state in which electrons bind (e&eb) magnetic flux lines to make (eb) new quasiparticles, and excitations have (e) a fractional elementary charge and possibly also fractional statistics. The 1998 Nobel Prize in Physics was awarded to Robert Laughlin, Horst Störmer, and Daniel Tsui for the discovery and explanation of the fractional Hall effect.[1] |
| (713) However, Laughlin's explanation was (=) a phenomenological guess and only applies to (e&eb) fillings {\displaystyle \nu =1/m} \nu =1/m where (e) {\displaystyle m} m is an odd integer. The microscopic origin of the FQHE remains unknown and is a major research topic in condensed matter physics. (Wikipedia) |
| (714)The quantum Hall effect (or integer quantum Hall effect) is a quantum-mechanical version of (e) the Hall effect, observed in (eb) two-dimensional electron systems subjected to (e&eb) low temperatures and strong magnetic fields, in which the Hall conductance σ undergoes (e&eb) certain quantum Hall transitions to take on (e&eb) the quantized values {\displaystyle \sigma ={\frac {I_{\text{channel}}}{V_{\text{Hall}}}}=\nu \;{\frac {e^{2}}{h}},} \sigma = \frac{I_\text{channel}}{V_\text{Hall}} = \nu \; \frac{e^2}{h}, where {\displaystyle I_{\text{channel}}} I_\text{channel} is the channel current, {\displaystyle V_{\text{Hall}}} V_\text{Hall} is the Hall voltage, e is the elementary charge and h is Planck's constant. The prefactor ν is known as the "filling factor", and can take on either integer (ν = 1, 2, 3, ...) or fractional (ν = 1/3, 2/5, 3/7, 2/3, 3/5, 1/5, 2/9, 3/13, 5/2, 12/5, ...) values. |
| (715) The quantum Hall Effect is referred to as (=) the integer or fractional quantum Hall effect depending on whether ν is an integer or fraction, respectively. |
| (716)The striking feature of the integer quantum Hall Effect is (=) the persistence of the quantization (i.e. the Hall plateau) as (e) the electron density is varied. |
| (717) Since the latter remains constant when (e) the Fermi level is in a clean spectral gap, it must be that the Fermi level is immersed in (eb) essential spectrum (i.e. of finite density of states), though Anderson localized. |
| (718) The quantization and stability of the Hall conductance in such conditions was achieved in 1994 by (e) Bellissard, van Elst and Schulz-Baldes using the methods of (e) Non-Commutative Geometry.[1] |
| (719)The fractional quantum Hall effect is more complicated, as its existence relies fundamentally on (e) electron–electron interactions. |
| (720)Although the microscopic origins of the fractional quantum Hall effect are unknown, there are several phenomenological approaches that provide (eb) accurate approximations. |
| (721) For example the effect can be thought of as (=) an integer quantum Hall effect, not of electrons but of (e) charge-flux composites known as (=) composite fermions. |

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| (722) In 1988, it was proposed that there was quantum Hall effect without Landau levels. This quantum Hall Effect is referred to as (=) the quantum anomalous Hall (QAH) effect. |
| (723)There is also a new concept of the quantum spin Hall effect which is an analogue of the quantum Hall effect, where (e) spin currents flow instead of (e&eb) charge currents.[2] |
| (724)In two dimensions, when classical electrons are subjected to (e&eb) a magnetic field they follow (eb) circular cyclotron orbits. |
| (725)When the system is treated quantum mechanically (eb) these orbits are quantized. |
| (726)The energy levels of these quantized orbitals take on (e) discrete values |
| (727):{\displaystyle \textstyle E_{n}=\hbar \omega _{c}(n+1/2),} \textstyle E_{n}=\hbar \omega _{c}(n+1/2), where ωc = eB/m is the cyclotron frequency |
| (728) These orbitals are known as Landau levels, and at weak magnetic fields, their existence gives rise to (eb) many interesting "quantum oscillations" such as the Shubnikov–de Haas oscillations and the de Haas–van Alphen effect (which is often used to map the Fermi surface of metals). |
| (729) For strong magnetic fields, each Landau level is (=) highly degenerate (i.e. there are many single particle states which have the same energy En). |
| (730)Specifically, for a sample of area A, in magnetic field B, the degeneracy of each Landau level is (=) {\displaystyle \textstyle N=g_{s} BA/\phi _ {0},} \textstyle N=g_{s} BA/\phi _ {0}, where gs represents (eb) a factor of 2 for spin degeneracy, and ϕ0 = 2·10−15 Wb is the magnetic flux quantum. |
| (731) For sufficiently strong B-fields, each Landau level may have (e) so many states that all of the free electrons in the system sit in only a few Landau levels; it is in this regime where (e) one observes the quantum Hall Effect. |
| (732)The integers that appear in the Hall Effect are (=) examples of **topological quantum numbers**. |
| (733) They are known in mathematics as the first **Chern numbers** and are closely related to (e&eb) **Berry's phase.** |
| (734)A striking model of much interest in this context is the Azbel-Harper-Hofstadter model whose quantum phase diagram is (=) the Hofstadter butterfly shown in the figure. (deleted due to spatial constraints) |
| (735)The vertical axis is (=) the strength of the magnetic field and the horizontal axis is the chemical potential, which fixes the electron density. The colors represent the integer Hall conductance. Warm colors represent positive integers and cold colors negative integers. The phase diagram is fractal and has structure on all scales. In the figure there is an obvious self-similarity. |
| (736)Concerning physical mechanisms, impurities and/or particular states (e.g., edge currents) are important for (e) both the 'integer' and 'fractional' effects. |
| (737) In addition, Coulomb interaction is also essential in (eb) the fractional quantum Hall effect. |
| (738)The observed strong similarity between integer and (=) fractional quantum Hall effects is explained by (e) the tendency of electrons to form bound states with (e&eb) an even number of magnetic flux quanta, called composite fermions. |
| (739)Duality is the circumstance in which cranking up interactions results in (eb) another system we understand, usually one that can be described in terms of (e&eb) another set of weakly coupled particles. |
| (740)When such a dual representation is found, (eb) it gives a solution to the strong coupling problem. |
| (741) A further reason that theoretical physicists get excited about dualities is that they undermine (e) the very notion of an “elementary particle.” |
| (742)In terms of the original variables, the dual, weakly coupled particles are (=) generally solitons—large, heavy, classical objects made up of (e) many quanta. |

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| (743)The remarkable thing is that solitons can become (=) "elementary," in the sense that the model that describes the system can be reformulated with the solitons as (=) the basic constituents. |
| (744)An example of a duality, which we will use below, arises in (eb) Maxwell's theory of electromagnetism, where (e) the strength of the interactions is determined by (e) the electric charge e. |
| (745)The duality exchanges electric and (e&eb) magnetic point charges, and the interaction strength in the dual description is 1⁄e, which is small when (e) e is large. |
| (746)This idea has been generalized to (e) many other field theories, thanks, in particular, to tools of (e) supersymmetry [4]. |
| (747)Another classic example of duality involves (e&eb) vortices. |
| (748) A vortex is a localized object in the plane around which a phase field winds (e&eb). (A phase field takes values on the circle.) |
| (749)Dualities relating particles and vortices have (e) a rich history in particle physics and condensed-matter physics. |
| (750) The simplest version relates the particles in the normal state of a superfluid to (e&eb) the vortices in a dual superconductor [5]. |
| (751)This charge-vortex duality has been (=) a workhorse in the study of two-dimensional condensed-matter systems, such as (=) fractional Quantum Hall States, superconducting thin films, and beyond-Landau critical phenomena [6]. |
| (752)Vortices can be (=) fermions (as they frequently are, for instance, in type-II superconductors). |
| (753) A dual description of a system with (e&eb) fermionic vortices thus requires (e) a new ingredient, since (e) it must include (e) such a fermionic particle. |
| (754)The new work [1–3] builds on (e&eb) earlier theoretical work [7] to make (eb) an appealing proposal for (e) the required duality: the fermionic vortex of QED in (eb) 2+1 dimensions is (=) a free Dirac fermion. The argument for this proposal builds on our understanding of TIs. |
| (755)TIs are distinguished only by (e) their edge states. |
| (756) The most famous example is a (3+1)-dimensional time-reversal invariant TI, which can host boundary states whose band structure is defined by (e) a single Dirac cone. |
| (757)A single time-reversal-symmetric Dirac cone is not possible without (e) the bulk TI, so its appearance in experiments provides (eb) a signature of the bulk phase. |
| (758) But the Dirac cone is not (e) the only possible signature. |
| (759) Deformations of the boundary conditions can dramatically change (e&eb) the spectrum of the edge states, yet preserve some essential "TI-ness." This essential property is (=) called an anomaly in high-energy theory. |
| (760)Max Metlitski of the University of California, Santa Barbara, and Ashvin Vishwanath of the University of California, Berkeley [1], and independently, Chong Wang and Todadri Senthil at the Massachusetts Institute of Technology, Cambridge [2], use this ambiguity in the edge theory and the electric-magnetic duality of an auxiliary theory of bulk electromagnetism coupled to (e&eb) the TI to argue for the duality. |
| (761)When I said that (3+1)-dimensional QED had a duality that interchanged the electric and magnetic charges, one could have (e) complained that the electron is (=) a fermion, while (e) a magnetic monopole may not be. |
| (762)However, the authors observe that in a TI, a charge-two monopole is (=) also a fermion. |

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| (763) By analyzing the spectrum of electric and magnetic charges, they conclude that the duality exchanges the electron and (e&eb) the charge-two monopole and, moreover, exchanges a TI with (e&eb) time-reversal symmetry with (e&eb) a "chiral TI," which instead has (e) particle-hole symmetry |
| (764). But this is just what is required of (e) a symmetry that exchanges a fermionic particle with (e&eb) a fermionic vortex: time-reversal symmetry maps the electron to itself but maps the vortex to (e&eb) another vortex whose phase winds (e&eb) the other way (see Fig. 1). |
| (765)The exchange of particles and holes, in turn, preserves (eb) the vortex. |
| (766)The bulk duality implies (eb) a duality between edge theories. |
| (767)The freedom to deform the edge of the TI implies (eb) a weak form of equivalence: although the edge theories are (=) topologically the same (in terms of symmetry assignments and (e&eb) anomalies), they need not have (e) identical low-energy physics. |
| (768)One side of the duality is (=) a free, massless Dirac particle, but the long-distance physics of (2+1)-dimensional QED is (=) a long-standing problem, which is important in the study of spin liquids. |
| (769) Many gauge theories in two dimensions develop (eb) a gap in their energy spectrum because of (e) the proliferation of monopoles. |
| (770) In the realization proposed by Metlitski and Vishwanath [1] and by Wang and Senthil [2], this cannot happen because (e) the monopoles carry electric charge, and indeed the stronger notion of duality as (=) an equivalence of low-energy physics may hold. |
| (771)Evidence for the stronger relation comes from (e) David Mross and colleagues at the California Institute of Technology, Pasadena [3], who have derived a lattice version of the duality using (e) a clever but explicit change of variables. |
| (772)Their construction is particularly interesting because (e) it uses (e) a lattice of one-dimensional models which themselves can only occur at the edge of a TI. |
| (773)Authors explicitly derive the duality between a free electronic Dirac cone and (e&eb) quantum electrodynamics in (eb) (2+1) dimensions (QED3) with (e&eb) N=1 fermion flavors. |
| (774)The duality proceeds via (e&eb) an exact, nonlocal mapping from electrons to (e&eb) dual fermions with (e&eb) long-range interactions encoded by (e) an emergent gauge field. |
| (775)This mapping allows (eb) us to construct (eb) parent Hamiltonians for exotic topological-insulator surface phases, derive (eb) the particle-hole-symmetric field theory of (e) a half-filled Landau level, and nontrivially constrain (e) QED3 scaling dimensions. |
| (776)They similarly establish duality between bosonic topological insulator surfaces and (e&eb) N=2 QED3. |

**Black hole thermodynamics with generalized uncertainty principle Li Xiang, X. Q. Wen (Submitted on 6 Jan 2009 (v1), last revised 18 Mar 2009 (this version, v2))**

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| (777) In the standard viewpoint, the temperature of a stationary black hole is proportional to (=) its surface gravity, $T_H=\hbar\kappa/2\pi$. |
| (778)This is a semiclassical result and the quantum gravity effects are not (e) taken into consideration. |
| (779)This Letter explores a unified expression for (e) the black hole temperature in the sense of (e) a generalized uncertainty principle (GUP). |
| (780) Discussion involves a heuristic analysis of (e) a particle which is absorbed by (e) the black hole. |
| (781)Besides a class of static and spherically symmetric black holes, an axially symmetric Kerr-Newman black hole is considered. Different from the existing literature, authors suggest (eb) that the black hole's irreducible mass represent (eb) the characteristic size in (eb) the absorption process. |

(782)The information capacity of a remnant is also discussed by (e) Bousso's D-bound in de Sitter spacetime. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP 0910:046,2009 DOI: 10.1088/1126-6708/2009/10/046 Cite as: arXiv:0901.0603 [gr-qc] (or arXiv:0901.0603v2 [gr-qc] for this version)

IF we define

$$[a_7]_7 = [f_{7*}]_7 + [f_{7*}]_7 + [f_{7*}]\,P^*_7$$

$$[a_7]_7 = \left[f^{\,7}_{7\,*}\right] P^*_7 - \frac{h^7_*}{h_*} P^*(0)$$

$$e=[\,g^7_*]_7\,Q^*_7$$

$$\eta_7(t) = \left[h_*\right]_7 \beta_7 \left(\theta_7(t)\right) - \mu_7(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\,[a_7]M_7 - [a_7]_7 W_7 - [n(t)]M_7$$

$$w_7(t) = eM_7$$

Write

$$A=\begin{pmatrix} -[a_7]_7 - [a_7]_7 \\ e \qquad\quad 0 \end{pmatrix}$$

$$B=(t)=\begin{pmatrix} [\eta(t)]_7 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_7 \\ W_7 \end{pmatrix}$$

The solution of (32)is

$$X_7(t) = e^{A,t}[X]_7(0) + \int_0 e^{At}$$

## PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION

## ATTRITION MODELS

## MODULE EIGHT

Jia li studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal.

### 70. MODEL FORMULATION

**Model is applied to systems given in the following:**

**Researchers blur the line between classical and quantum physics by connecting chaos and entanglement July 12, 2016 by Sonia Fernandez**

(4162) Experimental link between quantum entanglement (left) and (e&eb) classical chaos (right) found using (e) a small quantum computer. Credit: University of California - Santa Barbara Using a small

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| quantum system consisting of three superconducting qubits, researchers at UC Santa Barbara and Google have uncovered a link between aspects of classical and (e&eb) quantum physics thought to be unrelated: classical chaos and (e&eb) quantum entanglement. |
| (4163) Their findings suggest (eb) that it would be possible to use (e) controllable quantum systems to investigate certain fundamental aspects of nature. |
| (4164) "It's kind of surprising because (e) chaos is this totally classical concept—there's (=) no idea of chaos in a quantum system," Charles Neill, a researcher in the UCSB Department of Physics and lead author of a paper that appears in Nature Physics. |
| (4165) "Similarly, there's no concept of entanglement within (eb) classical systems. And yet it turns out that chaos and (e&eb) entanglement are really very strongly and clearly related." |
| (4166) Initiated in the 15th century, classical physics generally examines and describes (eb) systems larger than atoms and molecules. |
| (4167) It consists of hundreds of years' worth of study including Newton's laws of motion, electrodynamics, relativity, thermodynamics as well as chaos theory—the field that studies the behavior of (e) highly sensitive and unpredictable systems. |
| (4168) One classic example of chaos theory is the weather, in which a relatively small change in one part of (e) the system is enough to foil predictions—and vacation plans—anywhere on the globe. |
| (4169) At smaller size and length scales in nature, however, such as those involving (e&eb) atoms and photons and their behaviors, classical physics falls short. |
| (4170) In the early 20th century quantum physics emerged, with its seemingly counterintuitive and sometimes controversial science, including (e) the notions of superposition (the theory that a particle can be located in (eb) several places at once) and entanglement (particles that are deeply linked behave as (=) such despite physical distance from one another). |
| (4171) And so began the continuing search for connections between the two fields. All systems are fundamentally quantum systems, according Neill, but the means of describing in a quantum sense the chaotic behavior of, (e) say, air molecules in an evacuated room, remains limited. |
| (4172) Imagine taking a balloon full of air molecules, somehow tagging them so you could see them and then releasing them (eb) into a room with no air molecules, noted co-author and UCSB/Google researcher Pedram Roushan. |
| (4173) One possible outcome is that the air molecules remain clumped together (e&eb) in a little cloud following the same trajectory around the room |
| (4174) And yet, he continued, as we can probably intuit, the molecules will more likely take off in (eb) a variety of velocities and directions, bouncing off walls and interacting with (e&eb) each other, resting after the room is (=) sufficiently saturated with them. |
| (4175) "The underlying physics is (=) chaos, essentially," he said. |
| (4176) The molecules coming to rest—at least on (eb) the macroscopic level—is the result of (e) thermalization, or of reaching equilibrium after they have achieved (eb) uniform saturation within the system. |
| (4177) But in the infinitesimal world of quantum physics, there is (=) still little to describe that behavior. |
| (4178) The mathematics of quantum mechanics, Roushan said, do not allow for (e) the chaos described by Newtonian laws of motion. |
| (4179) To investigate, the researchers devised an experiment using (e) three quantum bits, the basic computational units of (e) the quantum computer. |

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| (4180) Unlike classical computer bits, which utilize (e) a binary system of two possible states (e.g., zero/one), a qubit can also use (e) a superposition of both states (zero and one) as a single state. |
| (4181) Additionally, multiple qubits can entangle, or link so closely that their measurements will (=) automatically correlate. |
| (4182) By manipulating these qubits with electronic pulses, Neill caused them to interact, rotate and evolve (e&eb, e&eb, eb) in the quantum analog of (e) a highly sensitive classical system. |
| (4183) The result is a map of entanglement entropy of (e) a qubit that, over time, comes to strongly resemble that of (e) classical dynamics—the regions of entanglement in (eb) the quantum map resemble the regions of (e) chaos on the classical map. |
| (4184) The islands of low entanglement in (eb) the quantum map are located in the places of (e) low chaos on the classical map. |
| (4185) "There's a very clear connection between entanglement and (e&eb) chaos in these two pictures," said Neill. |
| (4186) "And, it turns out that thermalization is (=) the thing that connects chaos and (e&eb) entanglement. |
| (4187) It turns out that they are actually the driving forces behind (eb) thermalization. |
| (4188) "What we realize is that in almost any quantum system, including on quantum computers, if you just let it evolve and you start to study what happens as a function of time, (eb) it's going to thermalize," added Neill, referring to the quantum-level equilibration. "And this really ties together the intuition between classical thermalization and (e&eb) chaos and how it occurs (eb) in quantum systems that entangle." |
| (4189) The study's findings have fundamental implications for (e) quantum computing. |
| (4190) At the level of three qubits, the computation is relatively simple, said Roushan, but as researchers push to build increasingly sophisticated and powerful quantum computers that incorporate (e) more qubits to study highly complex problems that are beyond the ability of classical computing—such as those in the realms of machine learning, artificial intelligence, fluid dynamics or chemistry—a quantum processor optimized for such calculations will be a very powerful tool. |
| (4191) "It means we can study things that are completely impossible to study right now, once we get to bigger systems," said Neill.More information: Nature Physics, DOI: 10.1038/nphys3830 , http://www.nature.com/nphys/journal/vaop/ncurrent/full/nphys3830.html Journal reference: Nature Physics search and more info website Provided by: University of California - Santa Barbara |

**Title: coherent control of macroscopic quantum states in a single-Cooper-pair box Authors: Nakamura, Y.; Pashkin, Yu. A.; Tsai, J. S. Publication: Nature, Volume 398, Issue 6730, pp. 786-788 (1999). (Nature Homepage) Publication Date: 04/1999 Origin: NATURE DOI: 10.1038/19718 Bibliographic Code: 1999Natur.398...786N**

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| (4192) A nanometre-scale superconducting electrode connected to (e&eb) a reservoir via (e&eb) a Josephson junction constitutes (e) an artificial two-level electronic system: a single-Cooper-pair box. |
| (4193) The two levels consist of (e) charge states (differing by 2e, where (e) e is the electronic charge) that are coupled by (e) tunnelling of Cooper pairs through (e&eb) the junction. |
| (4194) Although the two-level system is macroscopic, containing (e) a large number of electrons, the two charge states can be (=) coherently superposed. |

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| (4195) | The Cooper-pair box has therefore been suggested as (=) a candidate for a quantum bit or `qubit'- the basic component of a quantum computer. |
| (4196) | Here authors report the observation of quantum oscillations in (eb) a single-Cooper-pair box. |
| (4197) | By applying a short voltage pulse via (e&eb) a gate electrode, we can control (e&eb) the coherent quantum state evolution: the pulse modifies (e&eb) the energies of (e) the two charge states non-adiabatically, bringing them into (e&eb) resonance. |
| (4198) | The resulting state-a superposition of (e) the two charge states-is detected by (e) a tunnelling current through (e&eb) a probe junction. |
| (4199) | Results demonstrate electrical coherent control of (e) a qubit in a solid-state electronic device. |

**Superconducting quantum computing From Wikipedia, the free encyclopedia**

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| (4200) | Superconducting quantum computing is a promising implementation of (e) quantum information technology that involves (e&eb) nanofabricated superconducting electrodes coupled through (e&eb) Josephson junctions. |
| (4201) | As in a superconducting electrode, the phase and (e&eb) the charge are conjugate variables. |
| (4202) | There exist three families of superconducting qubits, depending on (=) whether the charge, (e&eb) the phase, or neither of (e) the two is good quantum number. |
| (4203) | These are respectively termed charge qubits (e&eb) flux qubits, and (e&eb) hybrid qubits. |

**Theory**

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| (4204) | Unlike many other physical implementations of a qubit which involve (e&eb) exclusively two-level systems (such as nuclear spin and (e&eb) photon polarization), the integrated quantum circuit involved in (e&eb) a superconducting qubit is (=) a multi-level system, of which only the first two levels are used as (=) the computational basis. |
| (4205) | A basic requirement for such an implementation is that the energy levels are (=) not uniformly spaced, so that (e) photons of a particular frequency which cause (eb) transition between the 0 and 1 levels do not(e) cause transitions from the first level to (e&eb) the higher levels as well. |
| (4206) | For electronic signals to be carried from one part of the circuit to (e&eb) another without energy loss and (e&eb) hence decoherence, the system also needs to be (=) non-dissipative, i.e., the metallic parts involved should have (e) zero resistance. |
| (4207) | These circuits currently need to be (=) operated at very low temperatures so that, 1) superconductivity is (=) realized, and 2) thermal fluctuations do not (e) cause transitions between (e&eb) energy levels. |

**Quantum LC Circuit**

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| (4208) | The simplest non-dissipative quantum circuit consists simply of (e) an inductor L {\displaystyle L} L and (e&eb) capacitor C {\displaystyle C} C, with (e&eb) the metallic wires connecting them being superconducting. |
| (4209) | Here the flux Φ {\displaystyle \Phi } \Phi through (e&eb) the inductor and (e&eb) the charge Q {\displaystyle Q} Q in the capacitor are (=) canonically conjugate variables which obey (eb) the commutation relation [ Φ , Q ] = i ℏ {\displaystyle [\Phi ,Q]=i\hbar } [\Phi ,Q]=i\hbar |

**The Hamiltonian associated with this system is**

(4210) H = Φ 2 2 L + Q 2 2 C {\displaystyle H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}} H={\frac {\Phi ^{2}}{2L}}+{\frac {Q^{2}}{2C}}

(4211) This gives rise to the energy levels E = ℏ ω 0 ( n + 1 2 ) {\displaystyle E=\hbar \omega _{0}(n+{\frac {1}{2}})} E=\hbar \omega _{0}(n+{\frac {1}{2}}) Here ω 0 = 1 L C {\displaystyle \omega _{0}={\frac {1}{\sqrt {LC}}}} \omega _{0}={\frac {1}{{\sqrt {LC}}}}

**Quantum gravity Carlo Rovelli (2008), Scholarpedia, 3(5):7117. doi:10.4249/scholarpedia.7117 revision #135539 [link to/cite this article] Post-publication activity Curator: Carlo Rovelli Dr. Carlo Rovelli, Center of Theoretical Physics, and Marseille, France**

(4212) Quantum Gravity is the name given to any theory that describes gravity in (eb) the regimes where quantum effects cannot be disregarded.

(4213) At present, there is no such a theory which is universally accepted and confirmed by experience. Therefore the term "Quantum Gravity" indicates (eb) more an open problem than a specific theory.

(4214) Several research lines, at different levels of development, offer tentative solutions to the problem. These tentative quantum-gravity theories are variously viewed as competing research directions, or as contributions to the common goal of finding the physically correct theory. The quest for the good quantum theory of gravity bears on a number of fundamental issues and it is sometimes presented as the most important open problem in fundamental physics: the "Holy Grail" of contemporary theoretical physics.

**The problem of quantum gravity**

(4215) The current theory of (e) gravity is general relativity, a theory developed by Albert Einstein at the beginning of the XX century, which has changed in depth our understanding of (e) what are space and time.

(4216) Similarly, quantum mechanics has changed in depth our understanding of matter, energy and causality. These two theories, today very well empirically supported, have opened (eb) a major conceptual revolution in physics.

(4217) However, they appear to be incompatible, at least at a first reading, because (e) each of the two is formulated on the basis of (e) principles that are explicitly contradicted by (e) the other theory.

(4218) Therefore in spite of the immense development of scientific knowledge it has delivered, the physics of the XX century has left us with (e&eb) a major confusion on what are the basic conceptual ingredients for (e) understanding the physical world.

(4219) In other words, the scientific revolution opened by general relativity and quantum mechanics at the beginning of (e) the XX century is not concluded yet, and a new synthesis is required. Quantum gravity, merging general relativity and quantum mechanics should be this new synthesis.

**Quantum space and quantum time**

(4220) Quantum gravity is expected to force us to further modifications of (e&eb) the concepts of space and time, along (e&eb) the direction opened by Einstein's general relativity, in order to make them fully compatible with quantum theory.

(4221) In general relativity, space and time lose (e) their properties of being a fixed framework in which the dynamical world is (=) immersed.

(4222) They are identified with (=) the gravitational field and acquire (e) dynamical properties: spacetime can "fold" and (e&eb) "stretch" like a rubber sheet.

(4223) When we take quantum mechanics into account, we realize that this "rubber sheet" is (=) in fact a quantum field, and therefore, like all quantum fields, it must have (e) a microscopic granular structure (like the photons forming the electromagnetic field) and a probabilistic dynamics.

(4224) Therefore quantum gravity is likely to be (=) the theory of granular and probabilistic "quantum space" and (e&eb) "quantum time". (Model is given in the next paper). Building the mathematical language and the conceptual structure for making sense of such notions of quantum space and quantum time is the challenge for a quantum theory of gravity.

**The Planck scale**

(4225) Simple dimensional arguments show that (eb) the physical phenomena where (e) quantum gravitational effects becomes relevant as those characterized by (e) the length scale lPlanck=ħG/c3−−−−−√∼10−33cm, called the "Planck length".

(4226) Here $\hbar$ is the Planck constant that governs (e&eb) the scale of the quantum effects, G is (=) the Newton constant that governs (e&eb) the strength of the gravitational force, and c is the speed of light, which governs (e&eb) the scale of the relativistic effects.

(4227) The Planck length is extremely small. To have an idea, the Planck length is (=) as many times smaller than an atom, as an atom is (=) smaller than the solar system.

(4228) Current technology is not yet capable of observing physical effects at (eb) scales that are so small (although several recent suggestions of how it could be possible to do so have appeared.)

(4229) Because of this, we have no direct experimental guidance for building a quantum theory of gravity. This is not by itself a complete impediment, because general relativity and quantum mechanics are themselves strong guidance for constructing the theory: several major advances in the history of physics have been obtained in (eb) the absence of new experiments, from (e) the effort of merging two empirically supported but apparently contradictory theories. (Examples are Newton's merge of Kepler's and Galileo's theories, Maxwell's merge of electric and (e&eb) magnetic theory, or Einstein's derivation of special relativity from (e) the apparent contradiction between electromagnetism and mechanics.) However, until genuine quantum gravitational phenomena are directly or indirectly observed, (eb) we cannot confirm or falsify any of the current tentative theories.

## Physical applications

Some phenomena where quantum gravity plays a major role are the following.

**The microscopic structure of spacetime**

(4230) If we could measure the geometry of space and time at the Planck scale, (eb) we should be able to see quantum gravitational effects.

(4231) Many arguments indicate (eb) that the Planck length may appear as a limit to the infinite divisibility of space, that is, as (=) a minimal length

(4232) Intuitively, any attempt to measure smaller distances would result in (eb) the concentration of (e) too much energy in too small a space, with the result of forming (eb) a micro-black-hole, effectively subtracting (e) the region from observation.

(4233) A minimal length would complete the turn of fundamental scales in Nature, together with (e&eb) the speed of light, which is the maximal velocity of a body, and the Planck constant, which is (=) the minimal amount of action exchanged between two systems.

**Early cosmology**

(4234) According to the currently standard cosmological model, the Universe was very dense and hot in the past. Extrapolating back the model, we encounter (e&eb) a singular point of infinite density, temperature and curvature conventionally denoted the Big Bang.

(4235) However, this final extrapolation is certainly incorrect because (e) quantum gravitational phenomena become (=) dominant when (e) the universe is very dense and hot, and these effects are not (e) included in the usual model.

(4236) A quantum theory of gravity is needed in order to take these effects into (e) account and study the early instants in the life of our universe.

(4237) Some current theories of gravity (in particular loop quantum gravity, see below), indicate (eb) that the singular Big-Bang point is never reached, and the current expansion of the universe might have been preceded by (e) a collapsing phase, and a "Big Bounce".

(4238) One of the major hopes of observing traces of quantum gravitational phenomena is (=) in this cosmological context, as traces of (e) early universe phenomena left in (eb) the cosmic background radiation currently under (e&eb) intense observation, or in the background gravitational wave radiation, which is expected to be observed in (eb) the next decade.

**Black holes**

(4239) Quantum gravity should play a role in (e&eb) several aspects of black-hole physics.

(4240) First, it should give a complete understanding of (e) the thermal radiation that black holes are expected to produce (eb) first computed by Stephen Hawking.

(4241) Second, Hawking's analysis shows that black holes carry (e&eb) enormous entropy about 10 to the power 77 for a solar mass black hole.

(4242) What is the statistical mechanical origin of this number which is enormous even by the standards of thermodynamics? Third, quantum gravity is expected to replace the infinite singularity that general relativity predicts at the center of black holes with (e&eb) a more physically reasonable picture.

(4243) Finally, the theory should explain what happens at the end of the Hawking evaporation of (e) a black hole.

**Astrophysical effects**

(4244) Several astrophysical quantum-gravitational effects have been suggested. None has been observed so far, but different calculations suggest (eb) that they might observed be in the near future.

(4245) An example is a small dependence of the speed of light on the color of the light, caused by (e) the Planck-scale granularity of space.

(4246) The effect is very small because of (e) the smallness of the Planck scale, but it might become detectable if (e) it is cumulated over a very long intergalactic path traveled by (e) the light. Observations for testing this prediction are ongoing.

**Main approaches**

(4247) The following approaches to quantum gravity open fascinating perspectives on the structure of Nature, and might be correct. But for the moment, none of these offers a complete theory of quantum gravity and none of these has been corroborated by experience. Because we do not have experimental data with direct bearing on quantum gravity, of necessity each program singles out certain aspects as being fundamental and focuses on them in the hope that other problems can be (=) easily sorted out later.

(4248) Opinions vary widely on the relative merit of (e) these approaches and the polemic, as often in science, is sometimes fiery.

(4249) The list that follows is roughly ordered according to (e&eb) the number of scientists working in each line of research. One should however remember that history teaches us that often a large majority of scientists has been working on the wrong tentative theory, with only a few working on the theory that ended up to be the correct: scientific truth is judged by experimental corroboration, not by a democratic counting of votes! These theories must all be considered tentative, at present.

**String theory**

(4250) String theory is not just a theory of quantum gravity in the strict sense, because (e) its objective is wider: the theory aims at giving a unified description of (e) the physical world, where all physical entities are understood as manifestations of (e) the motion of a single object: a string.

(4251) Gravity emerges in the theory as (=) one of the aspects of (e) the dynamics of the string.

(4252) String theory can be defined in terms of (e&eb) a perturbation expansion around (e&eb) a fixed spacetime.

(4253) Remarkably, in this formulation certain infinities that plague (e&eb) perturbative quantum general relativity do not (e) appear and there is an ongoing program to establish that each term in the series is (=) in fact finite.

(4254) However, when summed the entire series appears to be (=) divergent.

(4255) In any case, a definition of the theory as (=) a perturbation expansion is not sufficient for describing genuine quantum gravitational phenomena, which appear in (eb) the nonperturbative regime.

(4256) Therefore much of the research in this area in recent years focuses on (e&eb) nonperturbative effects.

(4257) Numerous indications point out to the existence of (e) a fundamental nonperturbative definition of (e) string theory, to which the various perturbative formulations should (eb) converge, and various partial attempts to define (eb) this nonperturbative theory are ongoing.

(4258) In this sought-for nonperturbative formulation, the characteristic features of quantum gravity become (=) manifest: for instance, the lack of a fixed background space and time, and (e&eb) the resulting conceptual difficulties. Research is currently active to try to find this fundamental description of string theory.

**Loop quantum gravity and spin foams**

(4259) Loop quantum gravity is the most developed attempt to solve (e) the specific problem of finding a quantum version of general relativity, and thus a genuine merge of the conceptual novelties represented by (e) general relativity and quantum mechanics.

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| (4260) The theory is consistent with the other fundamental physical theories (such as the standard model of particle physics), but it does not unify gravity and (e&eb) these theories in the sense of presenting them as manifestation of the dynamics of a single physical entity. |
| (4261) Loop quantum gravity is based in (eb) the formulation of general relativity developed by (e) Abhay Ashtekar, and offers a precise mathematical description of quantum spacetime. |
| (4262) The granular properties of space can be (=) explicitly computed. |
| (4263) In particular, a result of the theory is that the area and the volume of (e) any physical surface of region turns out to be (=) "quantized" (like the energy of a hydrogen atom), and the corresponding (e&eb) discrete values that area and volume can take have been computed within the theory. |
| (4264) The quantum states of physical space are described in the theory by (e) labelled graphs called spin networks. |
| (4265) Each node of a spin network represents (eb) an elementary "quantum of space", and the links between these indicates (eb) who is next to whom, building (eb) the spacial structure. |
| (4266) The main incompleteness of the theory regards the relation between the Planck scale and (e&eb) macroscopic physics, and the consistency of (e) its classical limit. |
| (4267) Related to loop quantum gravity is (=) a covariant approach to (e) quantum that goes under (e&eb) the name of spin foam formalism. |
| (4268) This is sometimes presented as (=) the path-integral or, covariant (Lagrangian) version of (e) the canonical (Hamiltonian) loop formalism. |

**Noncommutative geometry**

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| (4269) Einstein's discovery is that gravity, which is a dynamical field, is (=) the geometry of spacetime. |
| (4270) Quantum mechanics teaches us that dynamical quantities are (=) noncommutative --meaning that the outcome of the measurement of two of their properties depends on (eb) the order in which the measurement is performed. |
| (4271) It is therefore natural to suspect that the mathematics needed to describe quantum spacetime is (=) a noncommutative version of geometry. |
| (4272) A number of different formulations of such a noncommutative theory of geometry are under study. One line of research sees (eb) noncommutative spaces as (=) effective (approximate) descriptions of the deep quantum structure of spacetime and studies quantum field theory over (e&eb) such spaces. |
| (4273) Alain Connes has developed a beautiful mathematical framework for (e) describing noncommutative metric spaces. |
| (4274) This framework appears to shed surprising light on the physical geometry underlying (e&eb) the standard model of particle physics and its relation with (e&eb) general relativity, and opens an intriguing perspectives towards quantum gravity. |

**Others**

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| (4275) Among these, Roger Penrose has introduced an approach to quantum gravity denoted "Twistor theory", in which the primary elements for (e) the description of spacetime are not (e) the spacetime points but the light rays |

(4276) . Rafael Sorkin has introduced an approach denoted "Coset (or Causal Set) theory", where (e) quantum spacetime is described by (e) probabilistic superpositions of sets of points whose only structure is (=)their causal connection.

(4277) Finally, approaches that must be mentioned is (=) simply the possibility of taking general relativity and merge it with (e&eb) quantum mechanics following the conventional methods of quantum field theory, but circumventing (e) the traditional difficulties by (e) an intelligent retuning of (e&eb) the quantization method itself. Several authors have questioned the conventional assumption that this is manifestly impossible.

(4278) The common hope, today, is that the development of the theoretical research and, hopefully, some direct input from experiments and observations could soon bring to a convergence or a selection among these tentative theories, and therefore more clarity on the quantum nature of (e&eb) space and time.

Let $p_8(t, a_8)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and $P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(g_8)Q_8$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit creations which is assumed to be differentiable and satisfies

$$g_8(0) = 0, \quad g_8(P_8) > 0, \quad P_8 \text{ 0}, \qquad \lim_{P_8 \to \infty} g_8(P_8) = (K_8)_8 \text{ ∞}. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_8 + (\rho_a)_8 + D_8(a_8, P_8, Q_8)\rho_8 + 0, \qquad 0 < a_8 \text{ ∞},$$

$$\rho_8(t,0) = \int_0^\infty n_8(a_8, P_8, Q_8)\rho_8(t_8, a_8)da_8$$

Where $D_8$ is the mortality function and $n_8$ is the fertility function.

Put $_8$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_8}{dt}\right]_8 = x_8\left[\frac{dQ_8}{dt}\right]_8 - x_8\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_8}{dt}\right]_8 = x_3\left[\frac{dQ_3}{dt}\right]_8 - x_4\left[\frac{dQ_8}{dt}\right]_8$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_8}{dt}\right]_8 - x_6\left[\frac{dQ_3}{dt}\right]_8$$

Where $x_8, x_8, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_8$ and $g_8(P_8)Q_8$

Now when you add the term $(b_8Q_8)$ with $g_8(P_8)(Q_8)$ we are infect see is that $(b_8Q_8)$ is credited incrementally part of $g_8(P_8)(Q_8)$ or instant nearly, when subtracted from $\frac{dQ_8}{dt}$ , it means that $g_8(P_8)(Q_8)$ is slowly eaten by $\frac{dQ_8}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_8 = \mu_8(a_8) + (f_8)_8(P_8)_8 + (f_8)_8(Q)_8 \quad (2)$$

$$n_8 = \beta_8(a_8)h_8(Q_8), \quad (3)$$

Where $\mu_8$the age- dependent debit rate whit is $\mu_8(a_8)$ 0 $\forall a > 0$; $(f_8)_8$ is the density- dependent death rate with $f_8(0) = 0$, $(f_8^8)_8(P_8)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_8)_8(P_8) = (k_8)_8$ ; $f_8$ is the functional response of creditors to prey with $[f_8]_8(0) = 0, (f_8^8)_8(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_8)_8(Q_8) = [k_8]$ ; $\beta$ is the age

dependent birth rate with $\beta(a)$ 0, $\forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h^8(Q)$ 0, $\forall Q$ 0, and $\lim_{Q\to\infty} h(Q) = k_3$ 0.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_8 + (\rho_a)_8 = -\mu_8(a_8) + (f_8)_8 P_8) + \big(f_8)_8(Q_8)\big)(\rho_8)_8 \quad (4)$$

$$\rho_8(t_8, 0) = h_8(Q_8) \int_0^\infty \beta_8(a_8)\rho_8(t, a_8)da, \quad (5)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + g_8(P_8)Q_8 \quad (6)$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_8$ and $Q_8$ of (4-6) satisfy

$$0 \ \rho_8(t, a_8) \ \rho_8^*(t, a_8),$$

$$0 \ Q_8(t) \ Q_8^*(t),$$

Where $\rho_8^*(t, a_8)$ and $Q_8^*(t)$, are solutions of

$$(\rho_t^*)_8 + (\rho_a^*)_8 = -(\mu_8(a_8) + (f_8)_8(P_8^*) + (f_8)_8 Q_8^*))\,\rho_8^*, \quad (7)$$

$$\rho_8^*(t, 0) = \int_0^\infty \beta_8(a_8)\,\rho_8^*(t, a_8)da_8 \quad (8)$$

$$\frac{dQ_8^*}{dt} = -b_8 Q_8^* + g_8(P_8^*)Q_8^*, \quad (9)$$

With $P_8^*(t) = \int_0^\infty \rho_8^*\ (t, a_8)da_8$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_8^*(t)}{dt} = P_8^*(t)\big(r_8 - (f_8)(P_8^*) - (f_8)_8(Q^*)\big), \tag{10}$$

$$\frac{dQ_8^*(t)}{dt} = Q_8^*(t)(-b_8 + g_8(P_8^*)), \tag{11}$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_8(a)e^{-\int_0^a \mu_8(\tau)dt}e^{-ra}\ da_8 = 1. \tag{12}$$

If $r < 0$, the trivial solution $P_8^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(t)dt}\ da_8 > 1, \tag{$H1$}$$

Which guarantees $r_8 > 0$.

*Remark.* If $\beta_8(a) = [\beta_0]_8 a_8 e^{-\alpha_8 a_8}, \mu_8(a_8) \equiv \mu_8$ then the solution of (12) is

$$r_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_8]_8(P_8) = 0, [f_8](Q_8) = \lambda_8 Q_8\ h_8(Q_8) = 1, \qquad and\ \ g_8(P_8) = c_8 P_8$$

$$[\rho_8]_8 + [\rho_a]_8 = -(\mu_8 + \lambda_8 Q_8)\rho_8$$

$$\rho_8(t,0) = \int_0^{\infty} [\beta_0]_8 a_8 e^{-\alpha_8 a_8} \rho_8(t,a_8)\, da_8$$

$$\frac{dQ_8}{dt} = -b_8 Q_8 + c_8 P_8 Q_8$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_8}{dt} = \omega_8 P_8 - \lambda_8 P_8 Q_8$$

$$\frac{dQ_8}{dt} = b_8 Q_8 - c_8 P_8 Q_8$$

Where $\omega_8 = \sqrt{\beta_8} - \alpha_8 - \mu_8$

## 71. THE EQUILIBRIUM $(\rho_0)_8(a_8), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_8(a_8), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_8 + (\rho_a)_8 = -\big(\mu_8(a_8) + (f_8)_8(P_8)\big)\rho,$$

$$\rho_8(t,0) = \int_0^{\infty} \beta_8(a_8)\rho_8(t,a_8)da_8 \qquad (13)$$

The solution of (13) has the form of

$$\rho_8(t,a_8) = Ce_8^{r(t-a_8)} - \int_0^a (\tau)d\tau - \int_0^8 (f_8)_8\big(P_8(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_8(t) = C_8 e^{r_8 t - \int_0^t (f_8)_8\left(P_8(\tau)\right)d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau}\, da.$$

Hence, there exists an equilibrium $P_0 = \int_8^{-8}(r)$ if $k_8 > r$. This is equivalent to

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{(k_8)_8 a_8} da_8 < 1.$$

In fact, if we define

$$R_8(s_8) = \int_0^\infty \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau)d\tau} e^{-s_8 a_8}\, da_8$$

$R(s)$ is a decreasing function. $R_8(r_8) = 1$ And $R_8(k_8)_8 < 1$ imply $r_8 < (k_8)_8$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_8(a_8), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{k_8 a_8} da_8 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_8 \equiv \int_0^\infty (\rho_0)_{88}(a_8)da_8$ *satisfies*

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau)d\tau} e^{-(f_8)_8(P_0)a_8} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_8(a_8), 0)$

Let $u_8(t,a_8) = \rho(t,a_8) - \rho_0(a_8)$ Then

$$N_8(t) \equiv \int_0^\infty u_8(t,a_8)da_8 = \int_0^\infty \rho_8(t,a_8)da_8 - \int_0^\infty (\rho_0)_8(a_8)da = P_8(t) - P_0)_8$$

And

$$(u_8)_8 + (u_a)_8 = (\rho_t)_8 + (\rho_a)_8 - (\rho_0^8)_8(a_8)$$

$$= -\mu_8(a_8) + (f_8)_8(P_8) + (f_8)_8(Q_8))\,(u_8 + (\rho_0)_8) + \mu_8(a_8)(\rho_0)_8 + (f_8)_8(P_0)_8(\rho_0)_8$$

$$= -\mu_8(a_8)u_8 - (f_8)_8(P_8) + (f_8)_8(Q_8))\,u_8 - (\rho_0)_8\,(f_8)_8(P_8) - (f_8)_8(P_0)_8$$

$$+ (f_8)_8(Q_8)), \qquad (15)$$

Where $(\rho_0^8)_8(a_8)$ denotes $d(\rho_0)_8(a_8)/da_8$

Expansion $(f_8)_8(P_8)$ and $(f_8)_8(Q_8)$ at $(P_0)_8$ and 0 yield respectively

$$(f_8)_8(P_8) = (f_8)_8((P_0)_8 + N_8)$$

$$= (f_8)_8(P_0)_8 + (f_8^8)_8(P_0)_8N_8 + o_8(N_8), \qquad (16)$$

$$(f_8)_8(Q_8)$$

$$= (f_8^8)_8(0)Q + o_8(Q_8) \qquad (17)$$

Write $f_8(P_0) = f_{80}, f_8'(P_0) = f_{80}'$, and $f_8'(0) = f_{80}'$ and substitute (16) and (17) into (15) to obtain

$$(u_8)_8 + (u_a)_8 \approx -(\mu_8(a_8) + (f_{80})_8)u - (f_{80}^8)_8(\rho_0)_8(N_8) - (f_{80}^8)_8(\rho_0)_8(Q_8)$$

$$u_8(t,0) = \rho_8(t,0)(\rho_0)_8(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a) + (\rho_0)_8(a_8)da_8 - (\rho_0)_8\,(0)$$

$$= h_8(Q_8)\int_0^\infty \beta_8(a_8)(u_8(t,a_8)da_8 + (h_8(Q_8) - 1)(\rho_0)_8\,(0) \qquad (18)$$

Expanding

$$h_8(Q_8) = 1 + (h^8)_8(0)Q_8 + 0(Q_8) \equiv 1 + h_0^{'}Q_8 + 0Q_8$$

And substituting it into (18), we arrive at

$$u_8(t,0) \approx \int_0^{\infty} \beta_8\,(a_8)u_8(t,a_8)da_8 + (h_{a_8}^{'})(\rho_0)_8 Q_8$$

The linearization is the following :

$$[u_8] + [u_8] = -(\mu_8(a_8) + [t_{80}]_8 u_8 - [\alpha_8]_8(a_8)N_8 - [\alpha_8]_8(a_8)Q_8$$

$$u_8(t,0) = \int_0^{\infty} \beta_8(a_8)u_8(t,a_8)da_8 + [\xi_8]Q_8 \qquad (19)$$

$$\frac{dQ_8}{dt} = -b_8 Q_8$$

where

$$[\alpha_8]_8(a_8) = [t_8^8 a]_8[p,o]_8(a_8)d2(a_8) = [f_{80}^8]_8(\rho\, o\,)_8(a_8)\xi = [\mathrm{h\ o}^8]_8[\rho \mathrm{o}]_8(\mathrm{o}) \text{ and } \mathrm{b}_8 = \mathrm{b}_8 - g_8(P_8)$$

If $g_8(P_0) > b$, clearly, $(\rho_0)_8(a_8), 0)$ is unstable and the predator population is persistent. Assume $g_8(P_0)_8 < b$.

Then, $Q_8(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_8 + u_a = -(\mu_8(a_8) + (f_{80})_8)u_8 - (\alpha_8)_8(a_8)N_8$$

$$= -(\mu_8(a_8) + f_{80})u_8 - \alpha_8(a_8)\int_0^{\infty} u_8(t,a_8)da_8$$

$$u_8(t,0) = \int_0^\infty \beta_8\,(a_8)u_8(t,a_8)da_8$$

Using the method developed by Cushing [2, 3], we suppose

$$u_8(t,a_8) = \tilde{u}_8(a_8)^{e^{\gamma(t-a_8)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_8}\,da_8$. Then,

A solution of the first differential equation for $\tilde{u}(a_8)$ gives,

$$\tilde{u}_8(a_8) = e^{-\int_0^a \mu_8(\tau_8)d\tau - (f_{80})_8 a}\left(\tilde{u}_8(0) - w\int_0^a e^{\int_0^\tau \mu_8(s_8)ds + (f_{80})}\tau_8\,(\alpha_8)_8(\tau_8)e^{\gamma_8\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\right)\tilde{u}_8(0)$$

$$+w\int_0^\infty \beta_8(a_8)e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}\int_0^a (\alpha_8)_8(\tau_8)e^{-\gamma_8\tau_8}e^{(f_{80})_8\tau_8 + \int_0^\tau \mu_8(s_8)ds}d\tau_8\,da_8$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_8 a_8 - (f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau}da\right)(\tilde{u}_8(0))_8$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_8 a_8-(f_{80})_8 a_8-\int_0^a \mu_8(\tau_8)d\tau}\int_0^a e^{(f_{80})_8\tau_8+\int_0^\tau \mu_8(s_8)ds+y\tau}\alpha_8(\tau)d\tau\, da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_8 a_8-(f_{80})_8 a_8-\int_0^a \mu_8(t_8)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_8\,(a_8)Y_8(\gamma_8,a_8)da_8\right)u_8(0)$$

$$+w\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)\int_0^a(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8=0,$$

(24)

$$\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\tilde{u}_8(0)$$

$$+\left(-1-\int_0^{\infty}Y_8(\gamma_8,a_8)\int_0^a(\alpha_8)_8(\tau)_8(Y^{-8})_8(\gamma_8,\tau_8)d\tau_8\,da_8\right)w=0.$$

Equation (20) has a solution of the form of $u_8(t,a_8)=u_8(a_8)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_8(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_8(\gamma_8,a_8)=\int_0^a(\rho_0)_8(\tau_8)Y^{-8}\,(\gamma_8,a_8)d\tau_8$. We obtain the characteristic equation

$$f_{80_8}^{l}\left(\int_0^{\infty}Y_8(\gamma_8,a_8)da_8\right)\left(\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$+\left(1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8,a_8)da\right)\left(1+(f_{80}^{l})_8\int_0^{\infty}Y_8(\gamma_8,a_8)I_8(\gamma_8,a_8)da_8\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_8(a_8) = (\rho_0)_8(0)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - (f_{80})_8 a_8}$$

$$I_8(\gamma_8, a_8) = \int_0^a (\rho_0)_8(\tau_8)e^{\gamma_8\tau_8 + (f_{80})_8\tau + \int_0^\tau \mu_8(s_8)ds_8}d\tau$$

$$= \frac{(\rho_0)_8(0)}{\gamma_8}(e^{\gamma_8 a_8} - 1)$$

Let

$$E_8(a_8) = e^{-(f_{80})_8 a_8 - \int_0^a \mu_8(\tau_8)d\tau_8}$$

Then $[Y_8]_8[\gamma_8]_8, a_8) = E_8(a_8)e^{-\gamma_8 a_8}$. by a simple calculation,

$$\int_0^\infty \beta_8(a_8)Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma}\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}\ da\right);$$

$$\int_0^\infty Y_8(\gamma_8, a_8)I(\gamma_8, a_8)da = \frac{[\rho_0]_8(0)}{\gamma_8}\left(\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}\ da_8\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_8(0)}{\gamma_8}[f_{80}^8]\int_0^\infty E_8(a_8)e^{-\gamma_8 a_8}da_8\ (\int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8)$$

$$+\left(1 - \int_0^\infty \beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right).$$

$$\left(1+\frac{(\rho_0)_8(0)}{\gamma_8}(f_{80}^{'})_8\left(\int_0^{\infty}E_8(a_8)da_8-\int_0^{\infty}E_8(a_8)e^{-\gamma_8 a_8}da_8\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\right)(1+\frac{(f_{80}^{'})_8}{\gamma_8}(\rho_0)_8(0)\int_0^{\infty}E_8(a_8)da_8)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (24), to

$$1-\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)da_8=1-\int_0^{\infty}\beta_8(a_8)E_8(a_8)e^{-\gamma_8 a_8}da_8\neq 0,$$

Since $\int_0^{\infty}\beta_8(a_8)Y_8(\gamma_8 a_8)\int_0^{a}\alpha_8(\tau_8)Y^{-8}(\gamma_8,\tau_8)d\tau_8 da_8>0$. Therefore, the characteristic solution is

$$\gamma_8=-(f_{80}^{'})_8(\rho_0)_8\int_0^{\infty}E_8(a_8)da_8.$$

Since $f_8^{'}(P_0)_8>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_8),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_8),0)$ which is related to the extinction of predators and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $(p^*(a), Q^*)$. It must satisfy the equations:

$$[p_a^*]_8 + (\mu_8(a_8) + [f_8]_8(P_8^*) + [f_8]_8(Q_8^*))P_8^* = 0, \quad (26)$$

$$p_8^*(0) = h_8(Q_8^*)\int_0^\infty \beta_8(a_8)p_8^*(a_J)p_8^*(a_8)da_8 \quad (27)$$

$$g_8(P_8^*) - b_8 = 0, \quad (28)$$

Where $P_8^* = \int_0^\infty p_8^*(a_8)\, da_8$

Equation (28) yields $P^* = g^{-8}(b)$,

From (26),

$$p_8^*(a) = p_8^*(0)e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}$$

Substituting into (27), we need to solve

$$\left(1 - h_8(Q_8^*)\int_0^\infty \beta_8(a_8)e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(Q_8^*)a_8}da_8\right)p_8^*(0) = 0$$

For $Q_8^*$

Fix $p_8^*(0)$ and $P_8^*$ Let

$$H_8(z_8) = 1 - h_8(z_8)\int_0^\infty \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}da_{,8}$$

And assume that

$$[k_3]_8 \int_0^{\infty} \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(g_8^{-8}(b_8))a_8 - [k_8]_8 a_8}\, da_8 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty} [H_8(z_8)] > 0.$$

Clearly,

$$\frac{dH_8}{dz_8} = h_8(z)(\,[f_8^{'}]_8(z_8)\,) \int_0^{\infty} a_8 \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - f_8(P_8^*)a_8 - f_8(z_8)a_8} da_8$$

$-h^{'}(z_8) \int_0^{\infty} \beta_8(a_8)\, e^{\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8 - [f_8]_8(z_8)a_8}\; da_8$

By the assumptions on $\left[f_8^8\right]_8 and\ [h^8]^8 dH_8/dz > 0$. Hence, if

$$H_8(0) = 1 - \int_0^{\infty} \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8(P_8^*)a_8} da_8 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_8(0) < 0$ follows if $g_8^-(b_8) < [f_8^{-8}]_8(r_8)$, where $r_8$ is the solution of (12).

Indeed, since

$$R_8(s_8) = \int_0^{\infty} \beta_8(a_8)\, e^{-\int_0^a \mu_8(\tau_8)d\tau_8 e^{-y_8^a}} da_8$$

Is a decreasing function of $s_8$ when $P_8^* = [g^{-8}]_8(b_8) < [f_8^{-8}]_8(r_8)$, that is $[f_8]_8(P_8^*) < r$, $R\big(f_8(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_8(g^{-8}(b))) > 1, i.e.,$

$$\int_0^\infty \beta_8(a_8) e^{-\int_0^a \mu_8(\tau_8)d\tau_8 - [f_8]_8\left(g_8^{-8}(b)\right)a_8} > 1. \qquad \text{(H4)}$$

Then the system (4)-(6) has a unique positive equilibrium $(P_8^*(a_8), Q_8^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_8^*(a_8), Q_8^*)$ is more complicated. We again consider perturbations $m_8(t_8 a_8) = p_8(t_8, a_8) - p_8^*(a_8)$ and $W_8(t) = Q_8(t) - Q^*$.Then

$$M_8(t) \equiv \int_0^\infty m_8(t, a_8)da_8 = P_8(t) - P_8^*,$$

$$[m_8]_8 + [m_a]_8 = [p_8]_8 + [p_a]_8 - \left[p^*\right]_8(a_8) \approx -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8$$

$$-[f_{8*}^8]_8 p_8^*(a_8)M - [f_{8*}^8]_8 p_8^*(\mathrm{a})W_8,$$

$$m_8(t,0)p_8(t,0) - p_8^*(0) \approx h_* \int_0^\infty \beta_8(a_8)m_8(t,a_8)\,da + \frac{[h_*^8]_8}{[h_*]_8}\, p_8^*(0)W_8$$

$$\frac{dW_8}{dt} = \frac{dQ_8}{dt} \approx [g_*^8]_8 Q_8^* M_{,8}$$

Where $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $[f_{8*}]_8$ denoting $[f_8(P^*)]_8$, $\left[f_{8*}^{'}\right]_8$ denoting $\left[f^{'}{}_8(P^*)\right]_8\left[f_{8*}^{'}\right]_8$ denoting $[f_8]_8$ $(Q_8^*,$ $h_{*}^{'}$ denoting $h_*^{'}$ denoting $h_8^{'}$ $(Q_8^*)$, and $\left[g_*^{'}\right]_8$

Denoting $\left[\dot{g}\right]_8(P_8^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_8]_8 + [m_a]_8 = -(\mu_8(a_8) + [f_{8*}]_8 + [f_{8*}]_8)m_8 - \left[\dot{f}_{8*}\right]_8[p_8^*](a_8)M_8 - \left[\dot{f}_{8*}\right]_8 p_8^*(a_8)W_8$$

$$m_8(t,0) = [h_*]_8 \int_0^\infty \beta_8(a_8)m_8(t,a_8)da_8 + \frac{[h_*^8]_8}{[h^*]_8}[p^*]_8(0)\, W_{,8} \qquad (30)$$

$$\frac{dW_8}{dt} = [g_*^8]_8[Q_8^*]M_8$$

Integrating $[m_8]_8(t,a_8)$ with respect to a from 0 to ∞, we have

$$\frac{dM_8(t)}{dt} = \int_0^\infty [mt]_8(t,a_8)da_8$$

$$= -\int_0^\infty [m_a]_8\,(t,a_8)da_8 - \int_0^\infty \mu(a_8)m(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + \left[\dot{f}_{8*}\right]_8 \int_0^\infty p_8^*(a_8)da_8\right)M_8 - [f_{8*}^8]_8 \int_0^\infty p_8^*(a_8)da_8 W_8$$

$$= m_8(t,0) - \int_0^\infty \mu_8(a_8)\, m_8(t,a_8)da_8$$

$$-\left([f_{8*}]_8 + [f_{8*}]_8 + \left[\dot{f}_{8*}\right]_8 P_8^*\right)M_8 - \left[\dot{f}_{8*}\right]_8 P_8^* W_{,8}$$

From (30), we arrive at the equations

$$\frac{dW_8(t)}{dt} = \int_0^\infty ([h_*]_8\beta_8(a_8) - \mu_8(a_8))m_8(t,a_8)da_8$$

$$-\left([f_{8*}] + [f_{8*}]_8 + [f_{8*}']_8 P_8^*\right) M_8 + \left(\frac{[h_*']_8}{[h_*]^8}\, P_8^*(0) - f_{8*}' P_8^*\right) W_8,$$

(31)

$$\frac{dW_8(t)}{dt} [g'_*]_8\, Q_8^*\, M_{.8}$$

By the mean value theorem,

$$\int_0^\infty ([h_{"*}]_8 \beta_8(a_8) - \mu_8(a_8)) m_8(t, a_8) da_8 = ([h_*]_8 \beta_8(\theta_8) - \mu_8(\theta_8)) M_{,8}$$

Where $0< \theta_8 = \theta_8(t) < \infty$.

And if $a_8^8 [a_8^8] 2 < 4 [a_8] 2 < 4\, a_8 e$,

$\text{Re}\, \lambda_8 [\lambda_8]_8 = \text{Re} \lambda_8 [\lambda_8]_8 = -\frac{8}{8} a_8 [a2\ ]_8 < 0.$

Let $\Delta = \max_{8<u<\infty} |h_* [h_*]_8 \beta_8(a_8) - \mu_8(a_8)|$. Then $|n_8(t)|$ 24∀t 10. Hence we obtain the following result.

THEOREM 4. Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_8^8 < 4\, a_8 e$ and $[a_8] > 1\Delta$,or $[a_8^8]$ 24$[a_8]e$ and $(a_8 - \sqrt{(a_8^8)_8 - 4(a_8)_8 e) > 1\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+8}\right)$ and $\beta_8(a_8) = ae^{-ma}$ Then $\Delta = [\mu_0] + 2$ Simple calculation leads to

$$\int_0^{\infty} \beta_8(a_8)\mathrm{e}^{-\int_0^{\mathrm{a}} \mu_8(\tau_8)+d\tau_8}\mathrm{da} = \frac{1}{(\mathrm{m}+\mu_0+1)^8}+\frac{1}{(\mathrm{m}+\mu_0+1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_8 + [\mu_0]$ 0.5. Assume$[f_8](x_8)= K_8x_8$ Then the hypothesis (H2) is satisfied We take the conversion rate function $g_8(x_8)= [0x_8]/(2 + x_8)$; the functional responses $[[f_8](x) = [0]\, x_8/(1 + x_8)$and $h_8(x_8)$=$(1+\omega_8x_8)/(1 + x_8)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_8 + [\mu_0] + [\mathrm{K}][P_0])}{[m_8]_8 + (\mu_0)_8 + [\mathrm{K}][P_0])_8^3\ [\mathrm{K}]_8[P_0]_8^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_8$i e $[P_0] < b/(\delta_8 - b_8)$then the equilibrium$(P_0)$ $[P_0]_8(\mathrm{a}), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_8$

$[f_{8*}] = K_8\ \mathrm{P}^*\ [f_{8*}] = \frac{\sigma \mathrm{Q}_8^*}{8+\mathrm{Q}_8^*}$ $[h_*] = \frac{8+\omega_8\mathrm{Q}_8^*}{8+\mathrm{Q}_8^*}$

$[f_{8*}] = K_8\ [\mathrm{f}_{8*}^8] = \frac{\sigma_8}{(8+\mathrm{Q}_8^*)^8}[\ \mathrm{g}_*^8] = \frac{\sigma_8}{(8+\mathrm{P}_8^*)^8}$ $[\mathrm{h}_*^8] = \frac{\omega_8-8}{(8+\mathrm{Q}_8^*)^8}$

In order to have (H3) and (H4) hold,$m, \mu_0, \mathrm{K}, \mathrm{b}, \delta$, and $\omega$ need to satisfy

$$\frac{\omega_8}{(m+[\mu_0]_8+1+\frac{K_8b_8}{\delta_8+b_8}+\sigma)^8}+\frac{1\omega_8}{([m_8]_8+[\mu_0]_8+1+\frac{K_8b_8}{\delta_8+b_8}+\sigma_8)^3}<1$$

And

$$\frac{1}{(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^8} + \frac{1\omega_8}{(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8})^3} > 1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$(m_8 + [\mu_0]_8 + 1 + \frac{K_8 b_8}{\delta_8 - b_8} + (m_8 + [\mu_0]_8 + \frac{K_8 b_8}{\delta_8 - b_8} + \sigma_8)\, Q_8^*)^3$

$= m_8 + [\mu_0]_8 + 3 + \frac{K_8 b_8}{\delta_8 - b_8}(1 + Q^*)(1 + \omega_8 Q_8^*$

Now

$$\left[a_{11}\right] - 4\left[a_{8}\right] e = \left(\frac{1K_8 b_8}{\delta_8 - b_8}\right)^8 + \frac{4K_8 b_8 \sigma_8 Q_8^*}{(\delta_8 - b_8)(8 + Q_8^*)} + \frac{\sigma_8^8 Q_8^{*8}}{(1 + Q_8^*)^8}$$

$$\frac{4\, b_8 \sigma_8 (\delta_8 - b) Q^*}{\delta_8 (1 + Q_8^*)^8} - \frac{4\, (\delta_8 - b_8)^8 (1 - \omega_8) Q_8^*}{\delta_8 (1 + \omega_8 Q_8^*)(1 + Q_8^*)} P^*(0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_8\, ae_8^{mu}\ [\mu_0]_8 - \frac{a_8}{a_8 + 8}$$

Hence if $[a_8^8 > 4(a_8)_8 e$ and

$$[a_8] = \frac{1K_8 b_8}{\delta_8 - b_8} + \frac{\sigma_8 Q_8^*}{1 + Q_8^*} > 1\Delta$$

**The nontrivial equilibrium is stable.**

IF we define

$$[a_8]_8 = [f_{8*}]_8 + [f_{8*}]_8 + [f_{8*}] P_8^*$$

$$[a_8]_8 = \left[f_{8*}^{8}\right] P_8^* - \frac{h_*^8}{h_*} P^*(0)$$

$$e = [g_*^8]_8 Q_8^*$$

$$\eta_8(t) = \left[h_*\right]_8 \beta_8 \left(\theta_8(t)\right) - \mu_8(\theta(t)),$$

The system(32)can be written as the following system:

$$[M(t)]\left[a_8\right] M_8 - [a_8]_8 W_8 - [n(t)] M_8$$

$$w_8(t) = eM_8$$

Write

$$A = \begin{pmatrix} -[a_8]_8 - [a_8]_8 \\ e \qquad 0 \end{pmatrix}$$

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| $B=(t)=\begin{pmatrix} [\eta(t)]_8 & 0 \\ 0 & 0 \end{pmatrix}$ |
| And |
| $X(t)=\begin{pmatrix} M_8 \\ W_8 \end{pmatrix}$ |
| The solution of (32)is |
| $X_8(t) = e^{A,t}[X]_8(0) + \int_0 e^{At}$ |
| This equation notwithstanding its importance in all the models stand out for its importance for the following models<br><br>to which it is applied: |
| **Explicit Derivation of Duality between a Free Dirac Cone and Quantum Electrodynamics in (2+1) Dimensions David F. Mross, Jason Alicea, and Olexei I. Motrunich Phys. Rev. Lett. 117, 016802 – Published 27 June 2016** |

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| (482)These developments solve outstanding theoretical puzzles about (e&eb) interacting TIs and (e&eb) their symmetric, gapped, topologically ordered surface states [1, 2], and they have inspired furious activity in the field. |
| (483)The "auxiliary" bulk gauge theory formulated in Refs. [1, 2] actually predicts (eb) a novel (3+1)-dimensional spin liquid with (e&eb) time-reversal symmetry [8], which is (=) a candidate description of certain pyrochlore magnets. |
| (484) Perhaps most interesting is the connection to the manifestly particle-hole symmetric theory of (e) a half-filled Landau level [7], the best-understood non-Fermi liquid state. |
| (485) Briefly, the successful theory of this state makes a distinction between half-filled and (e&eb) half-empty electronic bands. |
| (486)The associated particle-hole symmetry is (=) emergent and an exact local realization can only occur at (eb) the edge of a TI |
| (487) The main outcome is a new understanding of (e) the Dirac-like nature of the composite fermion excitation of (e) this system, which provides (eb) a new point of view on the phenomenology [7, 9]. |
| (488)This circle of ideas has also led to (eb) new insights [10] on two-dimensional superfluid-insulator transitions with (e&eb) nearby metallic phases, where (e) the fermions come alive. |

(489)More recently, consequences of (e) the duality have been checked mathematically and numerically and the fermionic vortex duality has been fit into (e&eb) the wider web of field-theory dualities [11]. This research is published in Physical Review Letters, Physical Review B, and Physical Review X.

**Particle-vortex duality of two-dimensional Dirac fermion from electric-magnetic duality of three-dimensional topological insulators Max A. Metlitski and Ashvin Vishwanath Phys. Rev. B 93, 245151 – Published 27 June 2016**

(490)Particle-vortex duality is (=) a powerful theoretical tool that has been used to (e) study bosonic systems.

(491)Here, authors propose an analogous duality for (e) Dirac fermions in 2+1 dimensions.

(492) The physics of a single Dirac cone is proposed to be described by (e) a dual theory, QED3, with again a single Dirac fermion but coupled to (e&eb) a gauge field.

(493)This duality is established by considering two alternate descriptions of (e) the three-dimensional topological insulator (TI) surface.

(494)The first description is the usual Dirac fermion surface state. The dual description is accessed via (e&eb) an electric-magnetic duality of (e) the bulk TI coupled to (e&eb) a gauge field, which maps it to (e&eb) a gauged chiral topological insulator.

(495)This alternate description ultimately leads to (eb) a new surface theory, QED3, which provides (eb) a simple description of otherwise intractable interacting electronic states.

(496) For example, an explicit derivation of (e) the T-Pfaffian state, a proposed surface topological order of (e) the TI, is obtained by (e) simply pair condensing the dual fermions.

(497)The roles of time-reversal and (e&eb) particle-hole symmetries are exchanged by (e&eb) the duality, which connects some of results to (e&eb) a recent conjecture by Son on particle-hole symmetric Quantum Hall States. Received 15 April 2016 DOI:http://dx.doi.org/10.1103/PhysRevB.93.245151 ©2016 American Physical Society

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(498)Figure 1: The work of Imbrie [2] provides (eb) the proof that quantum systems can enter a (e&eb) many-body localized (MBL) phase, in which they cannot thermally equilibrate under (e&eb) their own internal dynamics.

(499)We learn in our statistical mechanics classes that a system of many interacting degrees of (e) freedom isolated from (e) the external environment will, under (e&eb) its own internal dynamics, reach (eb) thermal equilibrium after (e) a sufficiently long time.

(500)When this is true—as is often the case—the system serves as (=) a "bath" for its own subsystems: Any small part of the full system sees (e) the rest of the system as (=) a reservoir with which it can exchange (e&eb) energy, particles, and quantum entanglement. This process is called **thermalization.**

(501)However, there are some systems of interacting quantum particles or spins that are believed not to be able to serve as (=) a bath for themselves.

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| (502)Such many-body localized (MBL) [1] systems remain localized in states that are close to (e) their initial conditions (see Fig. 1). |
| (503)Bringing them to thermal equilibrium requires (e) an external bath. |
| (504) But so far, there has been no rigorous proof of (e) the existence of **such many-body localization**. Now, theoretical work by John Imbrie at the University of Virginia, Charlottesville [2], provides such a proof—although relying on (e) one very reasonable assumption. |
| (505)The concept of **localization** was introduced by Philip Anderson in 1958 [3], when (e) he showed that interacting spins in (eb) lightly doped semiconductors may not be able to (e) thermally equilibrate without (e) any help from phonons (the material's lattice vibrations). |
| (506)The relationship between Anderson localization and (e&eb) thermalization received little attention for almost half a century. |
| (507)But more recently, this topic was turned into a trending research field by great progress seen in (eb) many-body atomic, molecular, and optical physics, and quantum information. |
| (508)Such developments have allowed researchers to engineer and investigate (e&eb) experimentally a wide variety of systems—from trapped cold atoms to (e&eb) spin ensembles—containing (e) many interacting quantum degrees of (e) freedom that are (=) well isolated from (e) their environment. |
| (509)For quantum particles or waves in (eb) a disordered medium that do not (e) interact with each other, Anderson localization is well understood. |
| (510) It has been experimentally observed for phonons [4], photons [5], and ultracold atoms [6]. In 2006, Denis Basko, Igor Aleiner, and Boris Altshuler investigated whether an Anderson-localized system at (eb) nonzero temperature would remain localized when (e) small interactions between nearby particles are (=) introduced [7]. |
| (511) Studying such systems with perturbation theory, they found (eb) that this is indeed true to (e) all orders of approximation, thus establishing (eb) the concept of the MBL phase as (=) a state of matter at (eb) nonzero temperature. |
| (512) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators. |
| (513)However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8]. |
| (514)Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems. |
| (515) Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL. |
| (516)As a consequence, the system—evolving under (e&eb) its own Hamiltonian—cannot (e) reach a non-MBL state, and thermalization is (=) blocked. |
| (517)His conclusion holds for a broad class of 1D system with (e&eb) short-range interactions between particles or spins on a lattice. |
| (518) An example that could be realized experimentally with ultracold atoms is (=) a Bose-Hubbard system: Bosonic atoms hop on (eb) a 1D lattice in which the potential energy varies randomly from site to (e) site, and the atoms interact when (e) they are on the same site. |

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| (519)This result gives us confidence that some 1D systems undergo (e&eb) a true many-body localization phase transition between a delocalized or thermal phase (in which the system is (=) a bath for itself) and (e&eb) the MBL phase (in which it fails to be (=) a bath for itself). |
| (520)This novel quantum phase transition is a topic of active research, but it remains poorly understood [1]. Recent experiments—one with ions trapped in (e&eb) 1D (but with longer-range interactions than (e) those considered by Imbrie) [9] and others with (e&eb) ultracold atoms with (e&eb) short-range interactions (but trapped in (e&eb) more than one dimension) [10]—explored such a possible phase transition. |
| (521)However, in these cases, which are not covered by (e) Imbrie's proof, it remains possible that many-body localization is (=) not a true phase but (=) only a metastable regime: |
| (522) These systems might be (=) almost localized but still become (=) a bath for themselves on extremely long time scales (too long to be detected in these experiments) because of (e) nonperturbative effects. |
| (523)It is important to note that Imbrie's proof is valid under (e&eb) an assumption that is actually very reasonable, which he calls (=) "limited level attraction": the hypothesis that the eigenenergies of (e) the system's Hamiltonian do not (e) accumulate too strongly near (eb) degenerate points (states with almost the same energies). |
| (524) This is a realistic scenario for (e) all known systems. |
| (525)On the one hand, in systems that thermalize, near-degenerate eigenstates show (eb) energy-level repulsion due to (e) quantum-mechanical effects. |
| (526)On the other hand, in MBL systems, energy levels neither attract nor repel e*(e&eb) each other. |
| (527)No scenario with strong level attraction is known that would violate (e) Imbrie's assumption, but Imbrie could not mathematically prove that this seemingly obvious assumption is true. |
| (528) His proof thus remains, in principle, incomplete until this assumption is (=) proven or a general proof that does not (e) rely on this assumption is found. |
| (529)Many intriguing questions remain open. First, what is the nature of the quantum phase transition between quantum thermalization and (e&eb) many-body localization? |
| (530)Second, what are the possible systems in which there is (=) a sharply distinct MBL phase that is stable against (e&eb) all nonperturbative effects? |
| (531)Imbrie's proof partially answers the second question, saying that the MBL phase is indeed stable in (=) certain 1D systems with (e&eb) short-range interactions. |
| (532)So in those systems, the quantum phase transition does indeed exist. |
| (533)But Imbrie's proof does not work for (e) longer-range interactions or in more than one dimension. |
| (534)Does this limitation reflect (e&eb) the impossibility of a stable MBL phase outside of (e) the cases his approach can tackle, or is this (=) just a limitation of the mathematical approach he uses, and we may thus find stable many-body localization in (eb) a much wider class of systems? This research is published in Physical Review Letters. |

# PRESCIENTIAL DYNAMICS AND CONTENTIAL STATICITY OF ACCENTUATION ATTRITION MODELS

## MODULE NINE

**Abstract: Jiali studied predator – prey model in which the prey population is assumed to have age structure and is governed by the McKendrick- von Forester partial differential equation and the predator population is described by the classical Volterra- Lotka ordinary differential equation. Quite general hypotheses are assumed for the mortality function, the fertility function, and the functional responses of predation. Existence and stability of three biologically meaningful equilibria, corresponding to extinction of both species, persistence of one species prey, and coexistence of two species, are studied. A particular example and some numerical result are given. (Vide journal of mathematical analysis and applications 152, 399-415) 1990. We demonstrate its applicability cosmology/ quantum computation economics latal. Model is applied to following formulations Current bounds from (e) the polarization of the CMB predict (eb) the scale-invariant gravitational wave (GW) background of (e) inflation to be out of reach for (e) upcoming GW interferometers. This prospect dramatically changes if the inflaton is (=) a pseudoscalar, in which case its generic coupling to (e&eb) any abelian gauge field provides (eb) a new source of GWs, directly related to (e&eb) the dynamics of inflation. This opens up (eb) new ways of probing (e&eb) the scalar potential responsible for (e) cosmic inflation. Dividing inflation models into universality classes, authors analyze (e&eb) the possible observational signatures. One of the most promising scenarios is (=) Starobinsky inflation, which may lead to (eb) observational signatures both in (eb) direct GW detection as well as in (eb) upcoming CMB detectors. In this case, the complementarity between the CMB and (e&eb) direct GW detection, as well as the possibility of a multi-frequency analysis with (e&eb) upcoming ground and (e&eb) space based GW interferometers, may provide (eb) a first clue to the microphysics of (e) inflation. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph) Cite as: arXiv: 1603.01287 [astro-ph.CO] (or arXiv: 1603.01287v1 [astro-ph.CO] for this version Primordial gravitational waves for universality classes of pseudoscalar inflation Valerie Domcke, Mauro Pieroni, Pierre Binétruy (Submitted on 3 Mar 2016)STOP An analysis of general scalar-tensor gravitation theory, containing (e) two arbitrary functions of (e) the scalar field, is presented. The weak-field limit is considered in detail, and predictions for (e) the classical tests of gravitation theory are derived. A definite relationship between the light propagation and (e&eb) perihelion shift effects is found to hold under (e&eb) very general conditions. The theory of the detection of gravitational waves is also investigated, and the observable differences between the scalar and (e&eb) tensor components are indicated. Finally, the relationship between the properties of the source and (e&eb) its radiation is considered in the weak-source limit, and expressions for (e) the rate of energy loss are derived. It is shown that the existing observational data are consistent with (e&eb) the possibility that the scalar field represents (eb) a major component of gravitational radiation from (e) astronomical sources. Received 18 February 1970 DOI:http://dx.doi.org/10.1103/PhysRevD.1.3209 ©1970 American Physical Society Scalar-Tensor Theory and Gravitational Waves Robert V. Wagoner Phys. Rev. D 1, 3209 – Published 15 June 1970 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice (α=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society**

**Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP Masaru Shibata and Takashi Nakamura perform numerical simulations of (e) a three-dimensional (3D) time evolution of (e) pure gravitational waves. Authors use a conformally flat and K=0 initial condition for (e) the evolution of the spacetime. We adopt several slicing conditions to (e) check whether a long time integration is (=) possible in those conditions. For the case in which the amplitude of (e) the gravitational waves is low, (eb)a long time integration is possible by using (e) the harmonic slice and (e&eb) the maximal slice, while (e) in the geodesic slice ($\alpha$=1) it is not possible. As in the axisymmetric case and also in the 3D case, gravitational waves with (e&eb) a sufficiently high amplitude collapse by (e) their self-gravity and their final fates seem to be (=) as black holes. In this case, the singularity avoidance property of (e) the harmonic slice seems weak, so that (e) it may be inappropriate for (e) the formation problems of (e) the black hole. By means of the gauge-invariant wave extraction technique authors compute the waveform of (e) the gravitational waves at (eb) an outer region. They find that the nonlinearity of Einstein gravity induces (eb) the higher multipole modes even if only a quadrupole mode exists initially. Received 7 April 1995 DOI:http://dx.doi.org/10.1103/PhysRevD.52.5428 ©1995 American Physical Society Evolution of three-dimensional gravitational waves: Harmonic slicing case Masaru Shibata and Takashi Nakamura Phys. Rev. D 52, 5428 – Published 15 November 1995 STOP This paper is divided into four parts. In part A, some general considerations about (e&eb) gravitational radiation are followed by (e) a treatment of the scalar wave equation in the manner later to be applied to (e&eb) Einstein's field equations. In part B, a co-ordinate system is specified which is suitable for (e) investigation of outgoing gravitational waves from (e) an isolated axi-symmetric reflexion-symmetric system. The metric is expanded in (eb) negative powers of (e) a suitably defined radial co-ordinate r, and the vacuum field equations are investigated in detail. It is shown that the flow of information to infinity is controlled by (e&eb) a single function of (e) two variables called the news function. Together with initial conditions specified on a light cone, this function fully defines (eb) the behaviour of the system. No constraints of (e) any kind are encountered. In part C, the transformations leaving (eb) the metric in the chosen form are determined. An investigation of the corresponding transformations in (eb) Minkowski space suggests (eb) that no generality is lost by (e) assuming that the transformations, like the metric, may be expanded in (eb) negative powers of r. Gravitational Waves in General Relativity VII Waves from Axi-Symmetric Isolated Systems Royal Society Proceedings H. Bondi, M. G. J. van der Burg, A. W. K. Metzner Published 21 August 1962.DOI: 10.1098/rspa.1962.0161 STOP There is a deep tension between the well-developed theory of gravitational waves from (e) isolated systems and the presence of (e) a positive cosmological constant $\Lambda$, however tiny. In particular a generalization of Einstein's 1918 quadrupole formula that would allow (eb) a positive $\Lambda$ is not yet available. Authors first explain the principal difficulties and then show that it is possible to overcome (e) them in the weak field limit. These results also provide (eb) concrete hints for constructing the $\Lambda>0$ generalization of (e) the Bondi-Sachs framework for full, nonlinear general relativity. Received 18 October 2015 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.051101 © 2016 American Physical Society Gravitational Waves from Isolated Systems: Surprising Consequences of a Positive Cosmological Constant Abhay Ashtekar, Béatrice Bonga, and Aruna Kesavan Phys. Rev. Lett. 116, 051101 – Published 1 February 2016 STOP Searching for gravitational waves (GWs) from (e) binary black holes (BBHs) with LIGO and Virgo involves matched-filtering data against (e&eb) a set of representative signal waveforms—a template bank—chosen to (e) cover the full signal space of interest with (e&eb) as few template waveforms as possible. Although the component black holes may have (e) significant angular momenta (spin), previous searches for BBHs have filtered LIGO and Virgo data using (e) only waveforms where (e) both component spins are zero. This leads to (eb) a loss of signal-to-noise ratio for (e) signals where (e) this is not the case. Combining the best available template placement techniques and (e&eb) waveform models, authors construct a template bank of (e) GW signals from BBHs with component spins $\chi_{1,2}\in[-0.99,0.99]$ aligned (e&eb) with the orbital angular momentum, component masses $m_{1,2}\in[2,48]M_\odot$, and (e&eb) total mass $M_{total}\leq 50\ M_\odot$. Using effective-one-body waveforms with (e&eb) spin effects, authors show that less than 3% of (e) the maximum signal-to-noise ratio (SNR) of (e) these signals is lost due to (e) the discreetness of the bank, using (e) the early Advanced LIGO noise curve. Authors use simulated Advanced LIGO noise to (e) compare the sensitivity of this bank to (e) a nonspinning bank covering (e&eb) the same parameter space. In doing so, we consider the competing effects between improved SNR and (e&eb) signal-based vetoes and the increase in the rate of false alarms of (e) the aligned-spin bank due to (e) covering a larger parameter space. Authors find that the aligned-spin**

**bank can be (=) a factor of 1.3–5 more sensitive than a nonspinning bank to (e) BBHs with dimensionless spins >+0.6 and component masses ≳20 M⊙. Even larger gains are obtained for systems with (e&eb) equally high spins but smaller component masses. Implementing a search for gravitational waves from binary black holes with nonprecessing spin Collin Capano, Ian Harry, Stephen Privitera, and Alessandra Buonanno Phys. Rev. D 93, 124007 – Published 3 June 201**

## 72. MODEL FORMULATION

Let $p_9(t, a_9)$ denote the distribution function of debited types the number of debts with age from $\alpha$ to $\beta$ and

$P(t) = \int_0^\infty \rho(t,a)da$ is the total confederation size of debit summation. Denote the total number of by $Q$. Assume

that the credit population is governed by the classical Volterra-Lotka equation

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(g_9)Q_9$$

Where $b$ is the natural death rate in the absence of prey and $g(P)$ is the conversion rate function of debits into credit

creations which is assumed to be differentiable and satisfies

$$g_9(0) = 0, \quad g_9(P_9) > 0, \ P_9 \ 0, \qquad \lim_{P_9 \to \infty} g_9(P_9) = (K_9)_9 \ \infty. \qquad (1)$$

The debit confederation is assumed to be described by the McKendrick-von Forester equation

$$(\rho_t)_9 + (\rho_a)_9 + D_9(a_9, P_9, Q_9)\rho_9 + 0, \qquad 0 < a_9 \ \infty,$$

$$\rho_9(t,0) = \int_0^\infty n_9(a_9, P_9, Q_9)\rho_9(t_9, a_9)da_9$$

Where $D_9$ is the mortality function and $n_9$ is the fertility function.

Put $_9$ subscript to all parameters.

We assume following equations held

$$\left[\frac{dQ_9}{dt}\right]_9 = x_9\left[\frac{dQ_9}{dt}\right]_9 - x_9\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_9}{dt}\right]_9 = x_3\left[\frac{dQ_3}{dt}\right]_9 - x_4\left[\frac{dQ_9}{dt}\right]_9$$

$$\left[\frac{dQ_3}{dt}\right]_3 = x_5\left[\frac{dQ_9}{dt}\right]_9 - x_6\left[\frac{dQ_3}{dt}\right]_9$$

Where $x_9, x_9, x_3, x_4, x_5, x_6$ are accentuation and attention (sufficient credit- debit summation, radiation rate debits raised) , credit summation (amount of credits raised), the pathway that are naturally extinct has trivial and nontrivial equilibrium debit – credit co existence is also studied.

On similar liras we can its $bQ_9$ and $g_9(P_9)Q_9$

Now when you add the term $(b_9Q_9)$ with $g_9(P_9)(Q_9)$ we are infect see is that $(b_9Q_9)$ is credited incrementally part of $g_9(P_9)(Q_9)$ or instant nearly, when subtracted from $\frac{dQ_9}{dt}$ , it means that $g_9(P_9)(Q_9)$ is slowly eaten by $\frac{dQ_9}{dt}$ or instant nearly these factors are to be borne as we proceed to obtain conditions for co existence. When there is no existence of age structure, all the three categories in the nodule can be taken to be equivalent.

We assume that

$$D_9 = \mu_9(a_9) + (f_9)_9(P_9)_9 + (f_9)_9(Q)_9 \quad (2)$$

$$n_9 = \beta_9(a_9)h_9(Q_9), \quad (3)$$

Where $\mu_9$the age- dependent debit rate whit is $\mu_9(a_9)$ 0 $\forall a > 0$; $(f_9)_9$ is the density- dependent death rate with $f_9(0) = 0$, $(f_9^9)_9(P_9)0$, $\forall P$ 0, and $\lim_{p\to\infty}(f_9)_9(P_9) = (k_9)_9$ ; $f_9$ is the functional response of creditors to prey with $[f_9]_9(0) = 0, (f_9^9)_9(Q) > 0, \forall Q$ 0, and $\lim_{Q\to\infty}(f_9)_9(Q_9) = [k_9]$ ; $\beta$ is the age dependent birth

rate with $\beta(a) \geq 0, \forall a > 0$; and $h$ measures the effects of debit on the of credit with $h(0) = 1, h'(Q) \leq 0, \forall Q \geq 0$, and $\lim_{Q\to\infty} h(Q) = k_3 \geq 0$.

Summarizing these assumptions, we arrive at the following equations:

$$(\rho_t)_9 + (\rho_a)_9 = -\mu_9(a_9) + (f_9)_9 P_9) + \big(f_9)_9(Q_9)\big)(\rho_9)_9 \tag{4}$$

$$\rho_9(t_9, 0) = h_9(Q_9)\int_0^\infty \beta_9(a_9)\rho_9(t, a_9)da, \tag{5}$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + g_9(P_9)Q_9 \tag{6}$$

Using the comparison technique developed in Li [10], we can verify that the solution $\rho_9$ and $Q_9$ of (4-6) satisfy

$$0 \leq \rho_9(t, a_9) \leq \rho_9^*(t, a_9),$$

$$0 \leq Q_9(t) \leq Q_9^*(t),$$

Where $\rho_9^*(t, a_9)$ and $Q_9^*(t)$, are solutions of

$$(\rho_t^*)_9 + (\rho_a^*)_9 = -(\mu_9(a_9) + (f_9)_9(P_9^*) + (f_9)_9 Q_9^*))\,\rho_9^*, \tag{7}$$

$$\rho_9^*(t, 0) = \int_0^\infty \beta_9(a_9)\,\rho_9^*(t, a_9)da_9 \tag{8}$$

$$\frac{dQ_9^*}{dt} = -b_9 Q_9^* + g_9(P_9^*)Q_9^*, \tag{9}$$

With $P_9^*(t) = \int_0^\infty \rho_9^*\ (t, a_9)da_9$

From Eqs. (7)-(9), we can derive the following equation for *t* sufficiently large:

$$\frac{dP_9^*(t)}{dt} = P_9^*(t)\big(r_9 - (f_9)(P_9^*) - (f_9)_9(Q^*)\big), \quad (10)$$

$$\frac{dQ_9^*(t)}{dt} = Q_9^*(t)(-b_9 + g_9(P_9^*)), \quad (11)$$

Where *r* is the real solution of the equation

$$\int_0^\infty \beta_9(a)e^{-\int_0^a \mu_9(\tau)dt}e^{-ra}\ da_9 = 1. \quad (12)$$

If $r < 0$, the trivial solution $P_9^* = 0$ and $Q^* = 0$ is stable and, hence, two species go to extinction. This is trivial in mathematics. Therefore, we subsequently, throughout this paper, assume that

$$\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(t)dt}\ da_9 > 1, \quad (H1)$$

Which guarantees $r_9 > 0$.

*Remark.* If $\beta_9(a) = [\beta_0]_9 a_9 e^{-\alpha_9 a_9}, \mu_9(a_9) \equiv \mu_9$ then the solution of (12) is

$$r_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$$

Corresponding to the result in Gurtun and Levine [6]. In their model,

$$[f_9]_9(P_9) = 0, [f_9](Q_9) = \lambda_9 Q_9\ h_9(Q_9) = 1, \quad and\ \ g_9(P_9) = c_9 P_9$$

$$[\rho_9]_9 + [\rho_a]_9 = -(\mu_9 + \lambda_9 Q_9)\rho_9$$

$$\rho_9(t,0) = \int_0^{\infty} [\beta_0]_9 a_9 e^{-\alpha_9 a_9} \rho_9(t,a_9)\, da_9$$

$$\frac{dQ_9}{dt} = -b_9 Q_9 + c_9 P_9 Q_9$$

Which, obviously, is a special case of the model in this paper. They arrived at a pair of ordinary differential equations

$$\frac{dP_9}{dt} = \omega_9 P_9 - \lambda_9 P_9 Q_9$$

$$\frac{dQ_9}{dt} = b_9 Q_9 - c_9 P_9 Q_9$$

Where $\omega_9 = \sqrt{\beta_9} - \alpha_9 - \mu_9$

THE EQUILIBRIUM $(\rho_0)_9(a_9), 0)$

To study the case where only the debit confederation congregation the behavior of the equilibrium $(\rho_0)_9(a_9), 0)$ of (4)-(6).

Obviously, the existence of such equilibrium is equivalent to the existence of the equilibrium of the equations

$$(\rho_t)_9 + (\rho_a)_9 = -\big(\mu_9(a_9) + (f_9)_9(P_9)\big)\rho,$$

$$\rho_9(t,0) = \int_0^{\infty} \beta_9\,(a_9)\rho_9(t,a_9)da_9 \qquad (13)$$

The solution of (13) has the form of

$$\rho_9(t,a_9) = Ce_9^{r(t-a_9)} - \int_0^a (\tau)d\tau - \int_0^9 (f_9)_9\big(P_9(\tau)\big)d\tau$$

For $t$ sufficiently large, where $r$ is the solution of (12). By integrating with respect to a from 0 to ∞, WE OBTAIN

$$P_9(t) = C_9 e^{r_9 t - \int_0^t (f_9)_9(P_9(\tau))d\tau} \int_0^\infty e^{-ra}\, e^{-\int_0^a \mu(\tau)d\tau} da.$$

Hence, there exists an equilibrium $P_0 = \int_9^{-9}(r)$ if $k_9 > r$. This is equivalent to

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{(k_9)_9 a_9} da_9 < 1.$$

In fact, if we define

$$R_9(s_9) = \int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau)d\tau} e^{-s_9 a_9}\, da_9$$

$R(s)$ is a decreasing function. $R_9(r_9) = 1$ And $R_9(k_9)_9 < 1$ imply $r_9 < (k_9)_9$.

Thus, we can state an existence result of the equilibrium $(\rho_0)_9(a_9), 0)$ as follows.

THEOREM 1. *Assume that the hypotheses* (H1) *and*

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{k_9 a_9} da_9 < 1 \qquad (H2)$$

*Hold. Then the system* (4)-(6) *has the equilibrium*$(\rho_0(a), 0)$ *where*$(P_0)_9 \equiv \int_0^\infty (\rho_0)_{99}(a_9)da_9$ *satisfies*

$$\int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau)d\tau} e^{-(f_9)_9(P_0)a_9} da = 1 \qquad (14)$$

Next, we study the stability of $(\rho_0)_9(a_9), 0)$

Let $u_9(t,a_9) = \rho(t,a_9) - \rho_0(a_9)$ Then

$$N_9(t) \equiv \int_0^\infty u_9\,(t, a_9)da_9 = \int_0^\infty \rho_9(t, a_9)da_9 - \int_0^\infty (\rho_0)_9(a_9)da = P_9(t) - P_0)_9$$

And

$$(u_9)_9 + (u_a)_9 = (\rho_t)_9 + (\rho_a)_9 - (\rho_0^9)_9(a_9)$$

$$= -\mu_9(a_9) + (f_9)_9(P_9) + (f_9)_9(Q_9))\,(u_9 + (\rho_0)_9) + \mu_9(a_9)(\rho_0)_9 + (f_9)_9(P_0)_9(\rho_0)_9$$

$$= -\mu_9(a_9)u_9 - (f_9)_9(P_9) + (f_9)_9(Q_9))\,u_9 - (\rho_0)_9\,(f_9)_9(P_9) - (f_9)_9(P_0)_9$$

$$+ (f_9)_9(Q_9)), \qquad (15)$$

Where $(\rho_0^9)_9(a_9)$ denotes $d(\rho_0)_9(a_9)/da_9$

Expansion $(f_9)_9(P_9)$ and $(f_9)_9(Q_9)$ at $(P_0)_9$ and 0 yield respectively

$$(f_9)_9(P_9) = (f_9)_9((P_0)_9 + N_9)$$

$$= (f_9)_9(P_0)_9 + (f_9^9)_9(P_0)_9N_9 + o_9(N_9), \qquad (16)$$

$$(f_9)_9(Q_9)$$

$$= (f_9^9)_9(0)Q + o_9(Q_9) \qquad (17)$$

Write $f_9(P_0) = f_{90}, f_9'(P_0) = f_{90}'$, and $f_9'(0) = f_{90}'$ and substitute (16) and (17) into (15) to obtain

$$(u_9)_9 + (u_a)_9 \approx -(\mu_9(a_9) + (f_{90})_9)u - (f_{90}^9)_9(\rho_0)_9(N_9) - (f_{90}^9)_9(\rho_0)_9(Q_9)$$

$$u_9(t, 0) = \rho_9(t, 0)(\rho_0)_9(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t, a) + (\rho_0)_9(a_9)da_9 - (\rho_0)_9\,(0)$$

$$= h_9(Q_9)\int_0^\infty \beta_9(a_9)(u_9(t, a_9)da_9 + (h_9(Q_9) - 1)(\rho_0)_9\,(0) \qquad (18)$$

Expanding

$$h_9(Q_9) = 1 + (h^9)_9(0)Q_9 + 0(Q_9) \equiv 1 + h_0'Q_9 + 0Q_9$$

And substituting it into (18), we arrive at

$$u_9(t,0) \approx \int_0^\infty \beta_9(a_9)u_9(t,a_9)da_9 + (h_{a_9}')(\rho_0)_9Q_9$$

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PREDATOR-PREY POPULATION MODELS

The linearization is the following :

$$[u_9] + [u_9] = -(\mu_9(a_9) + [t_{90}]_9u_9 - [\alpha_9]_9(a_9)N_9 - [\alpha_9]_9(a_9)Q_9$$

$$u_9(t,0) = \int_0^\infty \beta_9(a_9)u_9(t,a_9)da_9 + [\xi_9]Q_9 \qquad (19)$$

$$\frac{dQ_9}{dt} = -b_9Q_9$$

where

$$[\alpha_9]_9(a_9) = [t_9^9a]_9[p,o]_9(a_9)d2(a_9) = [f_{90}^9]_9(\rho\, o\,)_9(a_9)\xi = [\mathrm{h\ o}^9]_9[\rho\mathrm{o}]_9(\mathrm{o}) \text{ and } \mathrm{b}_9 = \mathrm{b}_9 - g_9(P_9)$$

If $g_9(P_0) > b$, clearly, $(\rho_0)_9(a_9), 0)$ is unstable and the predator population is persistent. Assume $g_9(P_0)_9 < b$.

Then, $Q_9(t) \to \infty$. Thus, asymptotically the dynamics of (19) are governed by

$$(u_t)_9 + u_a = -(\mu_9(a_9) + (f_{90})_9)u_9 - (\alpha_9)_9(a_9)N_9$$

$$= -(\mu_9(a_9) + f_{90})u_9 - \alpha_9(a_9)\int_0^\infty u_9(t,a_9)da_9$$

$$u_9(t,0) = \int_0^\infty \beta_9\,(a_9)u_9(t,a_9)da_9$$

Using the method developed by Cushing [2, 3], we suppose

$$u_9(t,a_9) = \tilde{u}_9(a_9)^{e^{\gamma(t-a_9)}}$$

And define $w = \int_0^\infty \tilde{u}(a)e^{-\gamma a_9}\,da_9$. Then,

A solution of the first differential equation for $\tilde{u}(a_9)$ gives,

$$\tilde{u}_9(a_9) = e^{-\int_0^a \mu_9(\tau_9)d\tau-(f_{90})_9 a}\left(\tilde{u}_9(0) - w\int_0^a e^{\int_0^\tau \mu_9(s_9)ds+(f_{90})}\tau_9\,(\alpha_9)_9(\tau_9)e^{\gamma_9\tau}d\tau\right). \quad (21)$$

Substituting the expression (21) into the expression for above yields

$$\left(1 - \int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\right)\tilde{u}_9(0)$$

$$+w\int_0^\infty \beta_9(a_9)e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\int_0^a (\alpha_9)_9(\tau_9)e^{-\gamma_9\tau_9}e^{(f_{90})_9\tau_9+\int_0^\tau \mu_9(s_9)ds}d\tau_9\,da_9$$

$$= 0 \quad (22)$$

Substituting (21) into the definition of of *w* yields

$$\left(\int_0^\infty e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}da\right)(\tilde{u}_9(0))_9$$

$$+\left(-1-\int_0^{\infty} e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(\tau_9)d\tau}\int_0^a e^{(f_{90})_9\tau_9+\int_0^\tau \mu_9(s_9)ds+y\tau}\alpha_9(\tau)d\tau\,da\right)w=0.$$

(23)

Define $Y(\gamma,a)=e^{-\gamma_9 a_9-(f_{90})_9 a_9-\int_0^a \mu_9(t_9)d\tau}$. Then

$$\left(1-\int_0^{\infty}\beta_9\,(a_9)Y_9(\gamma_9,a_9)da_9\right)u_9(0)$$

$$+w\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)\int_0^a(\alpha_9)_9(\tau)_9(Y^{-9})_9(\gamma_9,\tau_9)d\tau_9\,da_9=0,$$

(24)

$$\left(\int_0^{\infty}Y_9(\gamma_9,a_9)da_9\right)\tilde{u}_9(0)$$

$$+\left(-1-\int_0^{\infty}Y_9(\gamma_9,a_9)\int_0^a(\alpha_9)_9(\tau)_9(Y^{-9})_9(\gamma_9,\tau_9)d\tau_9\,da_9\right)w=0.$$

Equation (20) has a solution of the form of $u_9(t,a_9)=u_9(a_9)e^{\gamma(t-a)}$ if and only if the system of equation (24) is solvable for $\tilde{u}_9(0)\not\equiv 0$ and $\gamma,w$ in the complwx plane.

Let $I_9(\gamma_9,a_9)=\int_0^a(\rho_0)_9(\tau_9)Y^{-9}\,(\gamma_9,a_9)d\tau_9$. We obtain the characteristic equation

$$f_{90_9}^{l}\left(\int_0^{\infty}Y_9(\gamma_9,a_9)da_9\right)\left(\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)I_9(\gamma_9,a_9)da_9\right)$$

$$+\left(1-\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9,a_9)da\right)\left(1+(f_{90}^{l})_9\int_0^{\infty}Y_9(\gamma_9,a_9)I_9(\gamma_9,a_9)da_9\right)$$

$$=0. \qquad (25)$$

Since

$$(\rho_0)_9(a_9) = (\rho_0)_9(0)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - (f_{90})_9 a_9}$$

$$I_9(\gamma_9, a_9) = \int_0^a (\rho_0)_9(\tau_9)e^{\gamma_9\tau_9 + (f_{90})_9\tau + \int_0^\tau \mu_9(s_9)ds_9}d\tau$$

$$= \frac{(\rho_0)_9(0)}{\gamma_9}(e^{\gamma_9 a_9} - 1)$$

Let

$$E_9(a_9) = e^{-(f_{90})_9 a_9 - \int_0^a \mu_9(\tau_9)d\tau_9}$$

Then $[Y_9]_9[\gamma_9]_9, a_9) = E_9(a_9)e^{-\gamma_9 a_9}$. by a simple calculation,

$$\int_0^\infty \beta_9(a_9)Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma}\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}\ da\right);$$

$$\int_0^\infty Y_9(\gamma_9, a_9)I(\gamma_9, a_9)da = \frac{[\rho_0]_9(0)}{\gamma_9}\left(\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}\ da_9\right).$$

Substituting in to (25), we get

$$\frac{[\rho_0]_9(0)}{\gamma_9}[f_{90}^9]\int_0^\infty E_9(a_9)e^{-\gamma_9 a_9}da_9\ (\int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9)$$

$$+\left(1 - \int_0^\infty \beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right).$$

$$\left(1+\frac{(\rho_0)_9(0)}{\gamma_9}(f_{90}^{'})_9\left(\int_0^{\infty}E_9(a_9)da_9-\int_0^{\infty}E_9(a_9)e^{-\gamma_9 a_9}da_9\right)\right)=0.$$

Hence the characteristic equation becomes

$$\left(1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\right)(1+\frac{(f_{90}^{'})_9}{\gamma_9}(\rho_0)_9(0)\int_0^{\infty}E_9(a_9)da_9)=0.$$

We need $\tilde{u}(a)\not\equiv 0$, namely $w\neq 0$. This necessarily leads, from (94), to

$$1-\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)da_9=1-\int_0^{\infty}\beta_9(a_9)E_9(a_9)e^{-\gamma_9 a_9}da_9\neq 0,$$

Since $\int_0^{\infty}\beta_9(a_9)Y_9(\gamma_9 a_9)\int_0^{a}\alpha_9(\tau_9)Y^{-9}(\gamma_9,\tau_9)d\tau_9 da_9>0$. Therefore, the characteristic solution is

$$\gamma_9=-(f_{90}^{'})_9(\rho_0)_9\int_0^{\infty}E_9(a_9)da_9.$$

Since $f_9^{'}(P_0)_9>0$, we obtain the following result

Theorem 2. Assume that the hypotheses (H1) and (H2) hold. Then the equilibrium $(\rho_0(a_9),0)$ is asymptotically stable if $g(P_0)<b$. It is unstable if $g(P_0)>b$.

## 3. COEXISTENCE

In the previous section, we studied the dynamic behavior of the trivial solution $(\rho_0(a_9),0)$ which is related to the extinction of predators (debit summation types: that which acts as a resource)and to

The possible persistence of the prey population. The purpose of this section is to investigate existence and stability of nontrivial equilibrium, which is related to coexistence of both credit summation and debit summation.

Denote the nontrivial equilibrium by $( p^*(a), Q^* )$. It must satisfy the equations:

$$[p_a^*]_9 + (\mu_9(a_9) + [f_9]_9( P_9^*) + [f_9]_9(Q_9^*) )P_9^* = 0, \quad (26)$$

$$p_9^*(0) = h_9(Q_9^*) \int_0^\infty \beta_9(a_9)p_9^*(a_J)p_9^*(a_9)da_9 \quad (27)$$

$$g_9(P_9^*) - b_9 = 0, \quad (28)$$

Where $P_9^*=\int_0^\infty p_9^*(a_9)\, da_9$

Equation (28) yields $P^* = g^{-9}(b)$,

From (26),

$$p_9^*(a) = p_9^*(0)e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}$$

Substituting into (27), we need to solve

$$\left(1 - h_9(Q_9^*)\int_0^\infty \beta_9(a_9)e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(Q_9^*)a_9}da_9\right)p_9^*(0) = 0$$

For $Q_9^*$

Fix $p_9^*(0)$ and $P_9^*$ Let

$$H_9(z_9) = 1 - h_9(z_9)\int_0^\infty \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9}\, da_{,9}$$

And assume that

$$[k_3]_9 \int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(g_9^{-9}(b_9))a_9 - [k_9]_9 a_9}\, da_9 < 1. \qquad (H3)$$

Then

$$\lim_{z\to\infty}[H_9(z_9)] > 0.$$

Clearly,

$$\frac{dH_9}{dz_9} = h_9(z)\left( [f_9']_9(z_9) \right)\int_0^\infty a_9\beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - f_9(P_9^*)a_9 - f_9(z_9)a_9} da_{\square}$$

$-h'(z_9)\int_0^\infty \beta_9(a_9)\, e^{\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9 - [f_9]_9(z_9)a_9}\ da_9$

y the assumptions on $\left[f_9^9\right]_9 and\ [h^9]^9 dH_9/dz > 0$. Hence, if

$$H_9(0) = 1 - \int_0^\infty \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9(P_9^*)a_9} da_9 < 0, \qquad (29)$$

There exists a unique positive solution $Q^*$

Now, we claim that $H_9(0) < 0$ follows if $g_9^-(b_9) < [f_9^{-9}]_9(r_9)$, where $r_9$ is the solution of (12).

Indeed, since

$$R_9(s_9) = \int_0^{\infty} \beta_9(a_9)\, e^{-\int_0^a \mu_9(\tau_9)d\tau_9 e^{-y\frac{a}{9}} da_9}$$

Is a decreasing function of $s_9$ when $P_9^* = [g^{-9}]_9(b_9) < [f_9^{-9}]_9(r_9)$, that is $[f_9]_9(P_9^*) < r$, $R\big(f_9(P^*)\big) > 1$ which implies $H(0) < 0$.

Summarizing the above statements, we have shown the existence of the nontrivial equilibrium.

THEOREM 3. Assume that the hypothesis (H3) holds and $R(f_9(g^{-9}(b))) > 1, i.e.,$

$\int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9)d\tau_9 - [f_9]_9\left(g_9^{-9}(b)\right)a_9} > 1.$ (H4)

Then the system (4)-(6) has a unique positive equilibrium $(P_9^*(a_9), Q_9^*)$.

Analyzing the stability of the nontrivial equilibrium $(P_9^*(a_9), Q_9^*)$ is more complicated. We again consider perturbations $m_9(t_9 a_9) = p_9(t_9, a_9) - p_9^*(a_9)$ and $W_9(t) = Q_9(t) - Q^*$.Then

$M_9(t) \equiv \int_0^{\infty} m_9(t, a_9) da_9 = P_9(t) - P_9^*,$

$$[m_9]_9 + [m_a]_9 = [p_9]_9 + [p_a]_9 - \left[p^*\right]_9(a_9) \approx -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9$$

$$-[f_{9*}^9]_9 p_9^*(a_9)M - [f_{9*}^9]_9 p_9^*(\mathrm{a})W_9,$$

$$m_9(t,0)p_9(t,0) - p_9^*(0) \approx h_* \int_0^{\infty} \beta_9(a_9) m_9(t, a_9)\, da + \frac{[h_*^9]_9}{[h_*]_9}\, p_9^*(0)W_9$$

$$\frac{dW_9}{dt} = \frac{dQ_9}{dt} \approx [g_*^9]_9 Q_9^* M_{,9}$$

Where $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $[f_{9*}]_9$ denoting $[f_9(P^*)]_9$, $[\dot{f}_{9*}]_9$ denoting $[\dot{f}_9(P^*)]_9[\dot{f}_{9*}]_9$ denoting $[f_9]_9$ $(Q_9^*$, $h_9^*$ denoting $\dot{h}_*$ denoting $\dot{h}_9$ $(Q_9^*)$, and $[\dot{g}_*]_9$

Denoting $[\dot{g}]_9(P_9^*)$.

Now, we study the asymptotic behavior of the solutions $m(t,a)$ and $W(t)$ of the system of equations:

$$[m_9]_9 + [m_a]_9 = -(\mu_9(a_9) + [f_{9*}]_9 + [f_{9*}]_9)m_9 - [\dot{f}_{9*}]_9[p_9^*](a_9)M_9 - [\dot{f}_{9*}]_9 p_9^*(a_9)W_9$$

$$m_9(t,0) = [h_*]_9 \int_0^\infty \beta_9(a_9)m_9(t,a_9)da_9 + \frac{[h_*^9]_9}{[h^*]_9}[p^*]_9(0)\,W_{,9} \qquad (30)$$

$$\frac{dW_9}{dt} = [g_*^9]_9[Q_9^*]M_9$$

Integrating $[m_9]_9(t,a_9)$ with respect to a from 0 to ∞, we have

$$\frac{dM_9(t)}{dt} = \int_0^\infty [mt]_9(t,a_9)da_9$$

$$= -\int_0^\infty [m_a]_9\,(t,a_9)da_9 - \int_0^\infty \mu(a_9)m(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + [\dot{f}_{9*}]_9 \int_0^\infty p_9^*(a_9)da_9\right)M_9 - [f_{9*}^9]_9 \int_0^\infty p_9^*(a_9)da_9 W_9$$

$$= m_9(t,0) - \int_0^\infty \mu_9(a_9)\, m_9(t,a_9)da_9$$

$$-\left([f_{9*}]_9 + [f_{9*}]_9 + [\dot{f}_{9*}]_9 P_9^*\right)M_9 - [\dot{f}_{9*}]_9 P_9^* W_{.9}$$

From (30), we arrive at the equations

$$\frac{dW_9(t)}{dt} = \int_0^{\infty} ([h_*]_9\beta_9(a_9) - \mu_9(a_9))m_9(t,a_9)da_9$$

$$-\left([f_{9*}] + [f_{9*}]_9 + \left[f_{9*}^{'}\right]_9 P_9^*\right) M_9 + \left(\frac{\left[h_*^{'}\right]_9}{[h_*]^9}\ P_9^*(0) - f_{9*}^{'} P_9^*\right) W_9,$$

(31)

$$\frac{dW_9(t)}{dt}\left[g^{'}{}_*\right]_9 Q_9^*\, M_{.9}$$

By the mean value theorem,

$$\int_0^{\infty} ([h_{"*}]_9\beta_9(a_9) - \mu_9(a_9))m_9(t,a_9)da_9 = ([h_*]_9\beta_9(\theta_9) - \mu_9(\theta_9))M_{,9}$$

Where $0< \ \theta_9 = \theta_9(t) < \infty$.

And if $a_9^9[a_9^9]2<4\ [a_9]2<4\ a_9e$,

$\text{Re}\,\lambda_9[\lambda_9]_9 = \text{Re}\lambda_9[\lambda_9]_9 = -\frac{9}{9}a_9[a2\ ]_9 < 0.$

Let $\Delta=\max_{9<u<\infty}\left|h_*[h_*]_9\beta_9(a_9) - \mu_9(a_9)\right|$. Then$|n_9(t)|$ ⍰⊿∀t ⍰0. Hence we obtain the following result.

**THEOREM 4**

Assume that the hypotheses(H3)and (H4)hold Then the nontrivial equilibrium$(\rho^*(a), \boldsymbol{Q}^*)$ is asymptotically stable if either $a_9^9 <4\ a_9e$ and $[a_9] > 1\Delta$,or $[a_9^9]$ ⍰4$[a_9]e$ and $(a_9 - \sqrt{(a_9^9)_9 - 4(a_9)_9e) > 1\Delta}$

These conditions sufficient The equilibrium might be unstable under some other circumstances in which cases Hopf-bifurcation might occur We intend to pursue this problem in the future.

Assume $\mu(a) = \mu_0 + \left(\frac{a}{a+9}\right)$ and $\beta_9(a_9) =$ ae-$^{ma}$ Then $\Delta$=$[\mu_0] + 2$ Simple calculation leads to

$$\int_0^{\infty} \beta_9(a_9) e^{-\int_0^a \mu_9(\tau_9) + d\tau_9 da} = \frac{1}{(m + \mu_0 + 1)^9} + \frac{1}{(m + \mu_0 + 1)^3}$$

Hence in order to satisfy the hypothesis (H2), we assume $m_9 + [\mu_0]$ 0.5. Assume$[f_9](x_9)= K_9 x_9$ Then the hypothesis (H9) is satisfied We take the conversion rate function $g_9(x_9)= [0x_9]/(2 + x_9)$; the functional responses $[[f_9](x) = [0]\, x_9/(1 + x_9)$and $h_9(x_9)$=$(1+\omega_9 x_9)/(1 + x_9)$(See Holling [7] and Rashevsky [23])

Equation (24),now has the form of

$$\frac{m_9 + [\mu_0] + [K][P_0])}{[m_9]_9 + (\mu_0)_9 + [K][P_0])_9^3\ [K]_9[P_0]_9^3} = 1$$

The equilibrium $[P_0]$ is the solution of this equation

If $\delta P_0/(1 + P_0) < b_9$i e $[P_0] < b/(\delta_9 - b_9)$then the equilibrium$(P_0)$ $[P_0]_9(a), 0)$ is stable This implies that the cedit confederation goes to extinction but that the cedit confederation persists $b_9$

$[f_{9*}] = K_9$ P$^*$ $[f_{9*}] = \frac{\sigma Q_9^*}{9+Q_9^*}$ $[h_*] = \frac{9+\omega_9 Q_9^*}{9+Q_9^*}$

$[f_{9*}] = K_9$ $\left[f_{9*}^{9}\right] = \frac{\sigma_9}{(9+Q_9^*)^9}$ $[g_*^9] = \frac{\sigma_9}{(9+P_9^*)^9}$ $\left[h_*^9\right] = \frac{\omega_9 - 9}{(9+Q_9^*)^9}$

In order to have (H3) and (H4) hold,$m, \mu_0$, K, b, $\delta$, and $\omega$ need to satisfy

$$\frac{\omega_9}{(m+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9+b_9}+\sigma)^9}+\frac{1\omega_9}{([m_9]_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9+b_9}+\sigma_9)^3}<1$$

And

$$\frac{1}{(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9})^9}+\frac{1\omega_9}{(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9})^3}>1$$

When the inequalities (33)and (34)hold there is a unique equilibrium $Q^*$ which can be solved from the following equation:

$$(m_9+[\mu_0]_9+1+\frac{K_9 b_9}{\delta_9-b_9}+(m_9+[\mu_0]_9+\frac{K_9 b_9}{\delta_9-b_9}+\sigma_9)\,Q_9^*)^3$$

$$=m_9+[\mu_0]_9+3+\frac{K_9 b_9}{\delta_9-b_9}(1+Q^*)(1+\omega_9 Q_9^*$$

Now

$$\begin{bmatrix}1\\1\end{bmatrix}-4\begin{bmatrix}a_9\end{bmatrix}e=\left(\frac{1K_9 b_9}{\delta_9-b_9}\right)^9+\frac{4K_9 b_9\sigma_9 Q_9^*}{(\delta_9-b_9)(9+Q_9^*)}+\frac{\sigma_9^9 Q_9^{*9}}{(1+Q_9^*)^9}$$

$$\frac{4\,b_9\sigma_9(\delta_9-b)Q^*}{\delta_9(1+Q_9^*)^9}-\frac{4\,(\delta_9-b_9)^9(1-\omega_9)Q_9^*}{\delta_9(1+\omega_9 Q_9^*)(1+Q_9^*)}P^*(0)$$

And

$$\Delta = \lim_{0<a<\infty} |[h_*]_9 \, ae^{mu}_9 \quad [\mu_0]_9 - \frac{a_9}{a_9+9}$$

Hence if $[a^9_9 > 4(a_9)_9 e$ and

$$[a_9] = \frac{1K_9 b_9}{\delta_9 - b_9} + \frac{\sigma_9 Q^*_9}{1 + Q^*_9} > 1\Delta$$

Some of the examples to which the model has been applied are:

**New look at the galaxy-shaping power of black holes**

(834) "It creates (eb) these bubbles, and the bubbles keep the plasma hot.

(835) That's what prevents galaxies from (e) becoming even bigger than they are now."

(836) Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space.

(837) In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene.

(838) Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve.

(839) "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara.

(840) "It's like nature's thermostat, which keeps these galaxies from (e) growing.

(841) If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars."

**A new look at the galaxy-shaping power of black holes**

(842) An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the

(843) McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said.

(844) Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.

**Doomed satellite missions**

(845) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.

(846) At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.

(847) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.

(848) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.

(849) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.

(850) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people

(851) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."

(852) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.

(853) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(854) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.

(855) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.

(856) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.

(857) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses.

| |
|---|
| (858) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well. |
| (859) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma. |
| (860) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6. |
| (861) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus. |
| (862) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (863) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

| |
|---|
| (864) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (865) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (866) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (867) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (868) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (869) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (870) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (871) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (872) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (873) We know dark matter exists only from (e) its gravitational effects. |
| (874) It holds together galaxies and (e&eb) galaxy clusters. |
| (875) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (876) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |

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| (877) | They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (878) | X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (879) | They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (880) | This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (881) | An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |
| (882) | Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (883) | "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (884) | The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (885) | Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (886) | Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (887) | "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (888) | Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (889) | "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (890) | Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) |

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| consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories. |
| (891) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen. |
| (892) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer. |
| (893) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants. |
| (894) These were treated as (=) "aware" and "unaware" trials, respectively. |
| (895) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it. |
| (896) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another. |
| (897) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion. |
| (898) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness. |
| (899) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said. |
| (900) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness." |
| (901) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner. |
| (902) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain. |
| (903) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said. |
| (904) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

(905) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter.

(906) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined.

(907) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions.

(908) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda.

(909) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general.

(910) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole.

(911) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes.

(912) Noda and his colleagues instead studied the change in emissions over (e&eb) time.

(913) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection.

(914) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1).

(915) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole.

(916) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved.

(917) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson’s problem**

(918) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking “palm tree” organization.

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| (919) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society |
| (920) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible? |
| (921) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem. |
| (922) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other. |
| (923) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible. |
| (924) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges. |
| (925) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century. |
| (926) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching. |
| (927) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom. |
| (928) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth. |

**Spheres on trees**

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| (929) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration. |
| (930) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement. |
| (931) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum. |
| (932) By visualizing the problem in this way, the researchers noticed that these particular graphs don’t (e) have lots of branches extending from lots of (e) other branches. |
| (933) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure. |

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| (934) The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network. |
| (935) As a result, it provides (eb) an important clue for finding the global minimum. |
| (936) It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum. |
| (937) This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large. |

**Thompson's problem 2**

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| (938) By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society |
| (939) The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima. |
| **(940)** The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node. |

**Implications for Thomson's and Smale's problems**

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| (941) Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before. |
| (942) "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around." |
| (943) As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem. |
| (944) "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view." |
| (945) In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges. |
| (946) From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear. |
| (947) "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK. |
| (948) "For larger numbers, there will likely be (=) multiple funnels. |
| (949) Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum." |

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| (950) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem. |
| (951) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said. |
| (952) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org |

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| The nontrivial equilibrium is stable. |
| IF we define $[a_9]_9 = [f_{9*}]_9 + [f_{9*}]_9 + [f_{9*}]\,P^*_9$ |
| $$[a_9]_9 = \left[f^{9}_{9*}\right]P^*_9 - \frac{h^9_*}{h_*}P^*(0)$$ |
| $$e=[g^9_*]_9\,Q^*_9$$ |
| $$\eta_9(t) = \left[h_*\right]_9\beta_9\left(\theta_9(t)\right) - \mu_9\left(\theta(t)\right),$$ |
| The system(32)can be written as the following system: |
| $$[M(t)]\left[a_9\right]M_9 - [a_9]_9W_9 - [n(t)]M_9$$ |
| $$w_9(t) = eM_9$$ |
| Write |
| $$A=\begin{pmatrix}-[a_9]_9 - [a_9]_9 \\ e \qquad 0\end{pmatrix}$$ |

$$B=(t)=\begin{pmatrix} [\eta(t)]_9 & 0 \\ 0 & 0 \end{pmatrix}$$

And

$$X(t)=\begin{pmatrix} M_9 \\ W_9 \end{pmatrix}$$

The solution of (32)is

$$X_9(t) = e^{A,t}[X]_9(0) + \int_0 e^{At}$$

**Note: There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and**

**sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit, facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by**

**term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation**

**alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation,**

**appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module**

**LEN NUNNEYS**

## Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

## Module numbered one

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems

and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \quad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1965); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be referred to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 6985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_6}{dt} = [N(t)]_6[(F_6)(R(t)]_6 - [M]_6[R(t)]_6$$

$$\frac{d[R(t)]_6}{dt} = [B]_6\,[(R)_6(t-T_6)] - [D]_6[(R(t)]_6 - [N(t)]_6\left[G\left(R(t)\right)_6\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_6}{dt} = a_6[N(t)]_6 - a_6[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_6 = a_3[N(t)]_3 - a_4[N(t)]_6$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_6 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

irst, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.6b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The

regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 6a is at the

**TIME DELAY**

FIG. 6. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.6$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**DELAY DUE TO PREDATOR MATURATION**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_6}{dt} = N_6(t - T_6)F_6(R_6(t - T_6) - N_6(t)M_6((R_6(t) \qquad (5)$$

$$\frac{dR_6(t)}{dt} = B_6\big(R_6(t) - D_6(R_6(t)\big) - N_6(t)G_6(R_6(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_6}{dR_6} + \frac{N_6 dG_6}{dR_6} - \frac{dB_6}{dR_6} > \frac{N_6 G_6}{2M_6}\left[\frac{dF_6}{dR_6} + \frac{dM_6}{dR_6}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_6/dR_6$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_6$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_6$) and the other derived from the non-delayed or "immediate" processes($I_6$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

ABSOLUTE STABILITY

$$L_6 = \frac{N_6 G_6}{2M_6}\frac{d(6nF_6)}{dR_6} (\geq 0) \qquad (6a)$$

$$I_6 = -\frac{N_6 G_6}{2M_6}\frac{d(6nM_6}{dR_6} (\geq 0) \qquad (6b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_6 > L_6 - I_6$, must be modified ahe $dM_6/dR_6$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_6 > L_6 - I_6, \quad \text{when} \quad P_6^2/4 > I \qquad (8a)$$

$$P_6 > L_6 - I_6 + \frac{(I_6 - P_6^2}{(L_6 - I)} \quad \text{when} \quad P_6^2/4 < I. \quad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

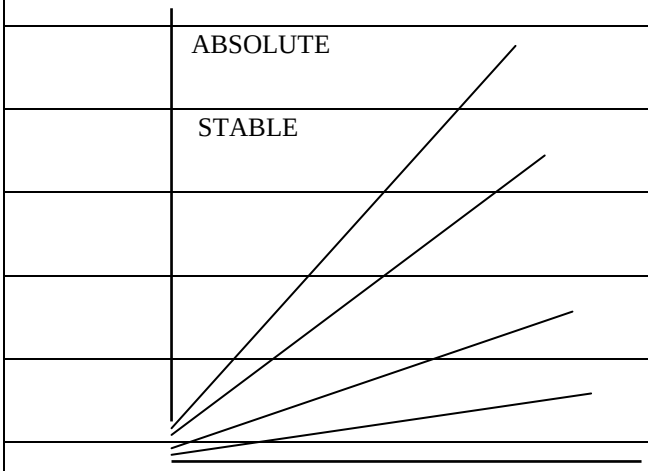



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

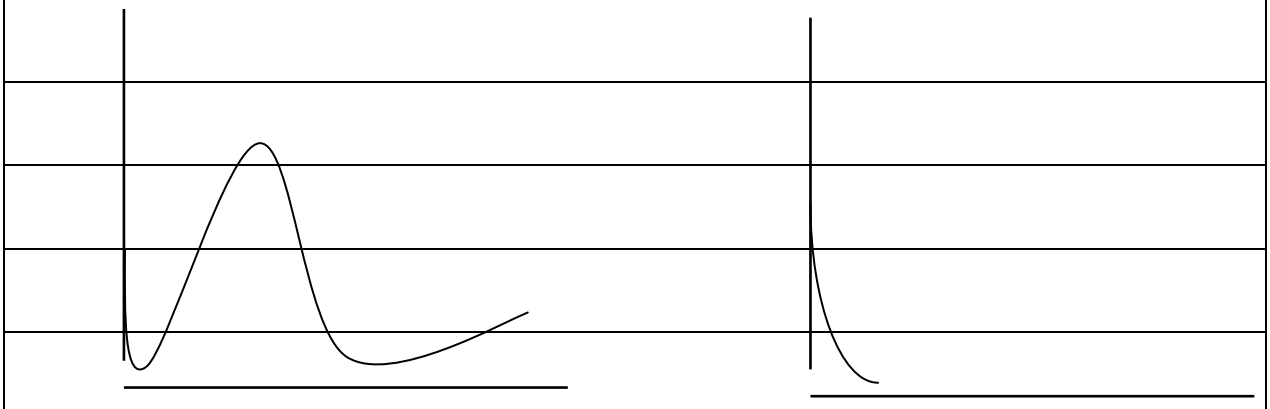
Can never completely offset the lagged effect,$L_6$; condition (8b) ensures that the total debit regulation, $P_6$, must always be positive even when $I_6$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_6$ and $L_6$. As $I_6$ increases relative to $L_6$ the zone of stability is enlarged; however,even when $I_6$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_6$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



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Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (6). Parameters in each case M=6; $dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6 = 0.05$ For (a) $NGdF/dR = 0.3, -N_6G_6dM_6/dR_6 = 0.6$; for (b)$NG_6dM_6/DR_6 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_6t}{dt} = N_6(t - T_6)F_6\big(R_6(t - T_6)\big) - N_6(t)M_6\big(R_6(t)\big)$$

(9)

$$\frac{dR_6t}{dt} = B_6\big(R_6(t - T_6)\big) - D_6\big(R_6(t)\big) - N_6(t)G_6\big(R_6(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) and the effective prey regulation ($dD_6/dR_6 + NdG_6/dR_6 - dB_6/dR_6$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we

could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

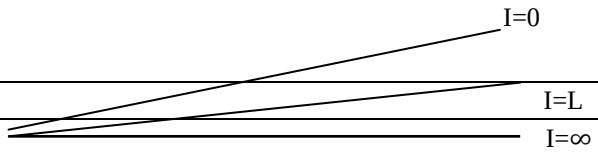
Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_6/dR_6$ is a measure of the control acting through the debit delay and setting $dB_6/dR_6 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_6/dR_6$ is a measure of the control acting through the credit delay although setting $dF_6/dR_6 = 0$ does not negate the effect of a predator delay; the value of $M_6$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_6/dR_6$and normalized by $M_6$ (see6a). The absolute stability of such a system is summarized in Fig. 6 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

TOTAL

ABSOLUTELY STABLE

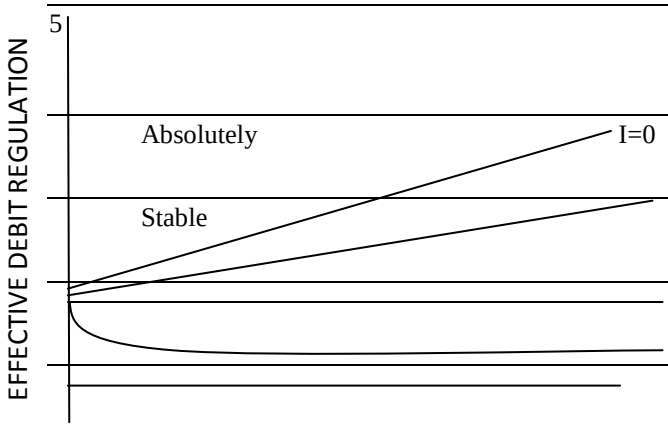



DELAYED DEBIT EFFECT

FIG6. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_6/dR_6 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.



-50 10

DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_6/dR_6 + N_6 dG_6/dR_6 + dB_6/dR_6)/M_6$). The immediate ($I_6$) and lagged ($L_6$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_6$(see 6b), is increased relative to the lagged part, $L_6$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_6 > I_6$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_6$is somewhat large than $I_6$ then the condition constraining prey regulation of recruitment ($dB_6/dR_6 < 0$) is

$$\frac{dD_6}{dR_6} + N_6\frac{dG_6}{dR_6} + \frac{dB_6}{dR_6} > \frac{(L_6 - I_6)M_6}{2} \qquad (10)$$

If $L_6 < I_6$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_6$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_6$

The - $D_6$ positions are defined by letting $\lambda_6 = iZ_6$ and separating (A1.2) into real and imaginary parts:

$$Z_6^2 + V_6 Z_6 \int_0^{\infty} \sin(Z_6 T_6) f(T_6) dT_6 - (X_6 + Y_6) = 0$$

$$Z_6 \left( U_6 - V_6 \int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6 \right) = 0$$

Noting that $X_6 + Y_6 = 0$ is we assume X+Y=0.It follows that $Z_6 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_6 = V_6 \int_0^{\infty} cos\left(Z_6 T_6\right) f\left(T_6\right) dT_6$$

If $U_6 > |V_6|$then given the constraints of $f(T_6)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A6.5)remains

greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

## ABSOLUTE STABILITY WITH A DELAY DUE TO CREDIT CONSOLIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_6^2 + (U_6 - V_6 + M_6)\lambda_6 - M_6\lambda_6 e^{-T_6\lambda_6} + [(U_6 - V_6)M_6 + Y_6] - [((U_6 - V_6 M_6 -]e^{-T_6\lambda_6} = 0$$

(For definitions see Appendix6) We follow the technique outlined in Appendix 6 and let $\lambda_6 = iZ_6$ and separate (A2.6) into real and imaginary parts

$$Z_6^2 + M_6 Z_6 \sin(Z_6 T_6) - (A_6 M_6 + Y_6) + (A_6 M_6 + X_6)\cos(Z_6 T_6) = 0$$

$$(A_6 M_6) Z_6 - M_6 Z_6 \cos(Z_6 T_6) + (A_6 M_6 - X_6)\sin(Z_6 T_6) = 0$$

Where $A_6 = U_6 - V_6$Using (A2.3), substitute for $Z_6^2$ and $Z_6$in (A2.2):

$$W_6^2(X_6^2 - 3M_6 A_6 X_6 - M_6^2(X_6 - Y_6) + 2M_6^3 A_6 + 2A_6^2 M_6^2]$$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_6 = 1 - \cos(Z_6 T_6)$.

In the absence of a delay the system is stable for $A_6 > 0,$ and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_6 > (X_6 - Y_6)/2\ M_6$ Thus a necessary condition for absolute stability is

$$A_6 > [\ X_6 - Y_6/2\ M_6$$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_6$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_6$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_6, A_6) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_6 \le 2$ *is* of interest First consider the extreme values

$$Q_6(0, A_6) = A_6^2\ (X_6 - Y_6)$$

$$Q_6(2, A_6) = (A_6 + 2\,M_6)^2\,(2A_6 M_6 - (X_6 - Y_6))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_6}{\partial A_6} = W_6^2\left(2M_6^3 + 4A_6 M_6^2 - 3M_6 X_6\right) + W_6\left(3M_6 A_6^2 - 2A_6 X_6 + 2M_6\left(2X_6 + Y_6\right)\right)$$

$$+2A_6\left(X_6+Y_6\right)$$

$=2M_6^2W_6^2(M_6+A_6)+M_6A_6^2W_6+\left(2A_6M_6-X_6+Y_6\right)(A_6W_6+M_6W_6^2$

$+(2-W_6)((2M_6X_6+M_6Y_6)W_6+A_6(X_6+Y_6)),$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_6$quadratic in $W_6$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_6$has no minimum in the range $0 \le W_6 \le 2$.Defining

$Q_6\left(W_6,A_6\right)=a_6W_6^2+c_6 where\ a_6\ b_6\ and\ c_6$ Are functions of $A_6$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_6 \ge b_6 \ge 0$$

And

$$b_6^2 \le 4a_6c_6$$

$$Case(i): Y_6=0$$

$Let\ Y_6 = 0$ And try the substitution A= $X_6 = /2M_6\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_6 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_6 X_6 /$ $2M_6$) is positive for *W<2. Since* $Q_6$ increases with $A_6$(see A2.9), increasing $A_6$will generate no further solutions of $Q_6$=0. Hence the condition A>X/2M defines absolute stability when $Y_6 = 0$.

$Case\ (ii)$: $Y_6 = y_6\ X_6$ ,

$Let\ Y_6 = y_6\ X_6$And try the substitution $A_6 = X_6/(2 + A_6\ y_6)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_6/(2 + n_6 y_6)$ we have

$$W_6^2\ y_6\left[X_6 n_6 \frac{(1 + n_6 y_6)}{(2 + n_6 y_6)} + M_6^2\left(2 - n_6 + n_6 y_6\right)\right]$$

$$-W_6\ X_6\left[\frac{X_6(1 + n_6 y_6)}{M_6^2(2 + n_6 y_6)^2} + 2\ y_6\left(n_6 - 1\right) + \frac{X_6^2(6 + y_6)}{M_6^2(2 + n_6 y_6)} = 0.\ (A2.12)\right.$$

Since $n_6$ and $A_6$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_6$

Defining $E_6 = (2 + n_6 y_6), the\ condition (A2.66?) can\ be\ written\ as$

$[\frac{X_6(E_6 - 6)}{E_6^2}]^2[\frac{X_6^2}{M_6^2} - 4\ y_6\ E_6^2\frac{X_6}{M_6^2} + 4y_6(n_6 - 2)E_6^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in ($X_6 / W_6^2$) can be expressed as

$2E_6[y_6E_6+(y_6E_6(6+y_6)(E_6-n_6))^{6/2}]>\frac{X_6}{M_6^2}>2E_6[y_6E_6-y_6E_6(6+y_6)(E_6-n_6))^{6/2}]$

**ABSOLUTE STABILITY**

Hence if $X_6/M_6^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_6$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_6-6)}{E_6^2}\left(\frac{X_6}{M_6^2}\right)^2-2y_6\frac{(E_6+n_6^2y_6)}{E_6}\left(\frac{X_6}{M_6^2}\right)-4y_6(E_6-n_6)$$ 0

And has the solution

$\frac{X_6}{M_6^2}>\frac{E_6}{(E_6-6)}[y_6(E_6+n_6^2y_6)+[y_6^2(E_6+n_6^2y_6)^2+4y_6(E_6-6)(E_6-n_6)]^{6/2}]$

Given $E_6>n_6$ this lower limit to $X_6/M_6^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_6/M_6^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_6$ in the appropriate range and this limit is further lowered by (A2.64)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_6<n_6$ $(ie.,2A_6M_6^2<X_6-Y_6)$ then (A2.6)and

(A2.8)show that $Q_6$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_6}{M_6^2} > 2E_6[y_6E_6 - (y_6E_6(6 + y_6)(E_6 - n_6)]^{6/2}]$$

This equations throws light on the individual general ledger and Nature's general ledger as we call it : see various papers on the issue. powers, which are the energies of Matrka we have already discussed, are not the only obstacles the yogi must overcome. He must, for example, also resist the temptation to rest content with the miraculous yogic powers (siddhi) he acquires in the course of his spiritual develop- ment. Again to do this he must practice Yoga. Similarly, in order to pervade the Fourth State gradually through the other states in the manner proper to practice at the Individual level, the method is the same. He must practise the higher yoga of the Tantras which, turning his mind inwards and freeing it from discursive representations, allows him to penetrate into the Supreme Principle. 312 Once the yogi has attained this contemplative state, his main problem is to make it permanent. In the introverted state the gross external movement of the breath is suspended and with it the activity of the intellect, mind, individualised consciousness, powers of the senses and the ego. 313 When the yogi rises out of this state, he is liable to fall again into the lower order of creation generated by Maya if he does not maintain his awareness of the higher reality he has experienced and allows his awakened, illumined insight to be obscured by the dream-like vision of thought-constructs. 314 Naturally, the yogi must rise out of the introverted condition of suspension. It is inherent in the very nature of reality that it should move out of itself. 315 Pure, universal consciousness initially transforms itself into the vital breath 316 charged with the impression iyasana) of the power of awareness attained through introversion. By attending to the pulse (spanda) of the breath as it moves out of the absolute, the

yogi can develop an intuitive sense of the inherent unity of all he will perceive in the mental and physical spheres created by the outpouring of consciousness. In this way he realises that his own nature is everywhere present in all he perceives and that all things thus reside within him. Blessed with this insight his consciousness remains free and unlimited even at the individual level where the breath, mind, senses and body are active. If the yogi fails to do this, he finds himself once again beset by the strictures of his embodied existence and must, as before, try to pervade all his other states of consciousness with the aesthetic delight (rasa) and wonder (camatkara) of the Fourth State he experienced in contemplation. Again this means that he must strengthen his pure, empowered awareness that his universal nature manifests as all things. 317 In this way he discovers Siva's presence in every sphere of individualised consciousness ranging from the breath to outer objectivity. The yogi's mind then becomes tranquil and undistracted because wherever it may wander, the yogi perceives only Siva, his authentic nature. 318 Consciousness is thus freed of all external referents and the yogi's subjectivity is purified of all identification with the body or anything else that belongs to the objective sphere. The yogi then becomes detached from the opposites of pleasure and pain and is transcendentally free (kevalin). 319 yogis are again, however, liable to fall if he allows himself to get entangled in the play of opposites. This fall is more serious than the others because, although he is caught by the confining restrictions of individual- ised consciousness as before, he is now also affected by karma. Fleeing from pain in the pursuit of pleasure he is bound to act (karma) to minimise one and maximise the other and so is thrown down to the lowest level of embodied subjectivity (sakala). In order to regain his lost state, he must ascend gradually, by Siva's grace, from one order of subjectivity to the next and so free himself progressively of the limitations of the lower levels to gain the greater freedom and expansion of the higher. As he progresses, the objective sphere also evolves from the grossest perceptions of physical objects outside the lowest order to subjectivity, through to the subtler inner, mental

perceptions to finally reach the order of subjectivity that contains objectivity within itself and is free to externalise it at will. 320 The degree to which this process develops depends, as before, on the yogi's awareness of the Fourth State. In consonance with the general principle that the remedy should suit the defect, the yogi is instructed to seek this higher state of consciousness in the wonder (camatkara) or delight (ananda) he feels in moments of intense physical pleasure. At first he experiences this subtle consciousness for an instant in the subjective sphere. If he manages to catch hold of it, it becomes more intense as the cognitive and objective spheres are also gradually pervaded and vitalised by it. Occasions for this practice are, for example, the sense of satisfaction one feels after a good meal or the aesthetic delight one experiences when listening to good music or the pleasure of sexual union with the Tantric consort or even solitary sexual excitation. In these moments of delight the yogi can penetrate momentarily into his own authentic Siva-nature (sambhavavesa) through the empowered contact (saktasparsa) m he makes with it in the freedom of the pure subjectivity of the Fourth State. 322 If the yogi develops his awareness of this higher level of consciousness and maintains it, he eventually experiences it constantly. 323 Clearly, what prevents the yogi from attending to his state of consciousness rather than the circumstances which induce it is the craving for pleasure (abhila§a) born of ignorance — the source of every impurity which clouds consciousness. Craving directs the yogi's attention towards outer, worldly things and so he is caught in the net of thought- constructs. 324 To free himself of his worldly desires and reverse this binding extroversion, the yogi must eradicate its cause. To be freed of all the ups and downs of the path and no longer be tormented by the possibility of a fall, the yogi must see reality perfectly and completely. This insight is itself liberation and the moment it dawns the yogi is instantly freed. This sudden realisation is the goal of Tantric Yoga. E Tantra declares: "He who perceives reality directly, even for the brief moment it takes to blink, is liberated that very instant and never reborn again." 325 Although the yogi's body and mind

continue to function as before, they are like mere outer coverings 326 which contain, but do not obscure, the mighty, universal consciousness which operates through them. The yogi's body is (=) the universe, the senses (=) the energies that (eb, eb+)vitalise it, his mind (=) Mantra, the rhythm of his breath (eb) the pulse of time and his inner nature pure, dynamic consciousness. Raised above all practice, and hence all possibility of falling to lower levels, the yogi realises that he has always been free 327 and that his journey through (e&eb) the **dark land of Maya was nothing but a dream, a construct of his own imagination** The channels of the senses by which the yogi enjoys the pleasure of the objects of sense are those through which he fills the three worlds with the consciousness emitted from his own Heart. 78 **The Doctrine of Vibration T**he senses are outer, subsidiary cycles of conscious energy (anucakra), revolving around the core or primary cycle (mukhyacakra) of pure consciousness. They scintillate with the diversity of sensations and resonate in the field of their awareness like the feathers of a peacock unfurled and folded back in the ecstasy of its dance. 79 Ultimately it is Siva Himself Who, for His own pleasure, 80 extends and withdraws them, 81 not the individual soul. The outward movement of the senses is the creative flow of consciousness from the subject who, like an ocean, is the source of the waves of their activity and the final resting place of the flow of the rivers of sensations and thought-forms that unfold through the arising of the power of each of the senses. 82 The Stanzas on Vibration declare: Indeed the individual soul does not activate the impulse of the will [which directs the body's activity] by himself alone but through his contact with [his] own [inner] strength {bald) made in such a way that he identifies with it [thus acquiring its power]. 83 The close contact between the senses and universal consciousness vitalises them and sets them in motion. They share in the properties of consciousness, just as a heated iron ball can burn, cook or give out light. 84 Spanda doctrines does not agree with the Sarnkhya view that the individual soul (purusa) alone stimulates activity in the objective sphere (prakrti) by the contact he makes with it. The soul can do

this only because universal consciousness animates him. Moreover, consciousness, which has the power to impel this activity, cannot itself be inactive as the Sarnkhya maintains. The influence it has on lifeless matter implies an active exertion. 85 This activity is the subtle, inner outpouring of the Self (svatmocchalatta) within the Self. Outside the confines of time and space, in the centre of the circle of the senses, this inner activity initiates the expansion (prasara) of the flow of conscious energy. As the pure exertion (udyoga, udyama) of Being (bhava), it is the source of the initial unfolding (unmesa) of Becoming. 86 Ksemaraja explains: [By] 'exertion' (udyama) is [meant] the emergence of the supreme [level of] intuition (pratibha) which is the sudden outpouring of the [immanent and transcendental] unfolding of the reflective awareness of pure consciousness. 87 The direct, intuition (pratibha) of reality, free of thought-constructs, is the supreme power (parasakti) of consciousness. The yogi can gain Divine **Body and Sacred Circle of the Senses 153** this pure intuitive awareness by grasping the initial expansion (unmesa) of consciousness which engenders the emergence of thought and the activity of the senses. It can be discovered in the subtle Centre, between one thought and the next. The Stanzas on Vibration teach: The expansion of consciousness that takes place when one is engaged in a single thought should be known to be the source from whence another arises. One should experience that for oneself! 88 This does not mean that the unfolding expansion of consciousness {unmesa) is an intermittent, transitory occurrence. 89 As the source of all thought, including the notions of past, present and future, 90 it stands outside the confines of time. Consciousness is in a state of perpetual expansion. It is an ever developing process, a tiny part of which we observe as the flow of events around us and the thoughts and feelings within us. The totality of this process can never be grasped objectively. Objectivity is a state of fragmentation (bheda) in which discrete elements are discernible and mutually distinguished as individual, specifically definable entities. It is the realm within which thought and language operate. To grasp reality in its completeness, we must go

beyond the partial representations of thought and speech. To experience the primordial source and basis of all things, we must pierce through the outer periphery of thought and plunge into the Centre, to discover the instant in which thought, and with it, the sense of diversity, initially emerges. Ksemaraja cites a practice from the Vijnanabhairava through which the yogi can achieve this state of awareness: Checking [the movement of] attention (cit) once it has quit its object, it must not [be allowed] to move to any other; contemplation (bhavana) then blossoms forth by [experiencing] the state in the centre between them (tanmadhyabhava). 91 Left unattended and undisciplined, awareness spontaneously shifts from one object of thought to (e&eb) the next. The transition from one to the other entails a movement through (e&eb) a state of pure indeterminate (nirvikalpa) awareness divested of all thought-forms and perceptions. The yogi who succeeds in checking the movement of his attention, experiences within himself the subsidence of thought into the emptiness of the Centre which unfolds full of the powerful pulsation of consciousness within which all objectivity merges and becomes one 92 in the trans- cendental outpouring (lokottarollasa) of its aesthetic rapture. 93 turning in on himself, the yogi must exert his power of awareness to (e) discern the vitality of (e) his own blissful Spanda nature. Just as a distant object becomes (=) more clearly visible by (e) exerting an effort to (e) perceive it, in the same way, through (e&eb) practice, the yogi's own Spanda nature becomes (=) progressively more evident to him, 94 The effort the yogi exerts to free himself of ignorance is essentially one with the force of awareness that impels the senses and mind. This inner exertion is the basis of every means to liberation. Ultimately bestowing the bliss of self- realisation, it carries the yogi along the path to salvation with all its twists and turns like wind which blows a cloth in various directions. 95 As the basis and ultimate goal of all yogic practice, this exertion is the pure wonder (camatkdra) inspired by the realisation of the essential identity of consciousness (cit) and being (sat). 96 As such it is the divine radiance (sphuratta) of the Heart of one's own consciousness 97 where the powers of

will, knowledge and action merge in the harmony (samarasya) of bliss. 98 Thus the highest form of worship (varivasya) is the contemplation of the inner strength (bald) or universal exertion of one's own true nature. 99 "One should examine [that reality]," says Ksemaraja, "with the exertion which is Bhairava's nature and is true service (sevana) to the integral inner nature [of all things] and the unfolding of the activity of one's own vitality (ojas) which withdraws all duality [back into consciousness]." 100 Mindful of the eternal joy (ahlada) of his inner strength, the yogi must give up all desire for the fleeting moments of petty pleasure (ksanikasukhalava) ]0] he may glean from outer objects. He must even give up the desire to achieve enlightenment and a reality (tattva) which can never be realised through a personal act of will. 102 In short, he must die to himself in an act of profound faith (sraddha) and adoration (adara) to find rest in the supreme reality free from the binding dichotomy of means and objectives (upeyopayabhava) — the reality of his perfected effort. The yogi, now full of devotion for this reality unfolding within himself and the universe (antarmukhatattva), merges into it and, by achieving the pure consciousness and bliss (cidananda) of the power (bala) of his own nature, achieves oneness. 103 The yogi is firmly established in his true nature by recognising the creative unfolding (unmesa) of consciousness through which he is able to know all things, whether in the present, past or future. 104 Master of the vitality which gives life to his mind and body, he can strengthenit directly without need of food or drink, 105 or quit it at will to vivify Divine Body and Sacred Circle of the Senses 155 that of another. 106 Experiencing the eternal satisfaction of his own uncreated nature, 107 the yogi is freed of disease and the infirmities of old age. 108 Inspired by this vitality, he has the inner strength to overcome the severest hardships and obstacles in the way of his ultimate goal. 109 The Stanzas on Vibration explain: Lassitude ravages the body and this arises from ignorance but if this is eliminated by an expansion of consciousness (unmesa), how can [ignorance], deprived of its cause, continue to exist? 110 Doubt (sanka) is the source of every spiritual ailment. 111 A

man in doubt about his true nature, and the path to follow by which he may come to realise it, is constantly overcome by difficulties. 112 "Due to ignorance" teaches the Sarvacaratantra, "the deluded man is in doubt and thus [suffers the cycle] of birth and death." 113 Caught between conflicting alternatives, the mind shifts from one to the other becoming more entangled in its own thoughts as it does so. Doubt thus contracts consciousness 114 sullying it with the turbulence of confused thinking. 115 At once both the root of "the ancient tree of transmigration" and the first sprout of its seed, 116 doubts deprives man of the innate bliss (sahajananda) of his own nature. 117 Like a thief it steals away the wealth of true knowledge and reduces man to an imagined state of poverty in which he feels hemmed in on all sides by constraints and limitations. 118 Ignorance and its sister, doubt, are the essence of all impurities (mala) which sully consciousness. To the degree in which the yogi is cured of this ailment, namely, the lassitude (glani) of doubt, his true nature becomes manifest, just as raw gold when heated is freed from dross. When the yogi recognises the all-powerful expansion (unmesa) of his consciousness, he is projected beyond the realm of relative distinctions between virtue and vice. Thus overcoming all doubt as to what he should or should not do, he penetrates into the pulsing Heart of Bhairava's consciousness. 119 As the Stanzas on Vibration say: An individual who [though] desirous of doing various things [but] is incapable of doing them due to his innate impurity, [experiences] the Supreme state when the disruption (ksobha) [of his ego] ceases. Then [the soul realises] that his [true] uncreated nature is [universal] agency and perceiving subjectivity and so he knows and does whatever [he] desires. 120 the individual soul is(=) a combination of (e) conscious (cit) and 156 The Doctrine of Vibration unconscious (acii) elements 121 brought together when he allows the power of his authentic nature to be obscured by (e) his own impurity. This impurity is (=) three-fold. The first is technically called (=) the 'Impurity of Individuality' (anavamald). Due to (e) this impurity, the individual soul fails to (e) recognise his all-embracing fullness (purnata) and is disturbed by (e) his craving

for (e) experience (bhoga) based on (e) the unreasoned assumption that he is (=) incomplete (apurnd) and in need of (e) something outside himself. 122 Thus it manifests as (=) desire or attachment (rdga). ]23 At the same time it is (=) primordial ignorance, 'the cause of the sprout of (e) transmigration'. 124 As such, it works in two ways. Firstly, it deprives (e) individual consciousness of freedom. The yogi affected by (e&eb) this impurity is incapable of (e) entering or emerging out of the states of contemplation as he wishes. Secondly, this impurity renders (eb) the individual unconscious of his freedom, which for(e) the yogi means that he is (=) liable to lose consciousness of (e) his true nature when he is (=) not in a state of (e) contemplative absorption. 125 The second impurity is (=) called the 'Impurity of Maya' (maylyamald). This comes into operation when consciousness has been contracted by (e) the Impurity of Individuality. It limits (e) the power of knowledge operating in (e&eb) the individual soul through (e&eb) the inner and outer organs of (e) perception. 126 Due to this impurity he perceives (e) diversity everywhere. 127 Finally, the third impurity, that of Karma, comes into play when (e) the individual, deprived of (e) the freedom and knowledge of (e) unconditioned consciousness, acts (e&eb) in his limited way prompted by (e) desires and fears for (e) his personal gain. Disturbed by (e) these three impurities the soul cannot find rest in (eb) himself.

And

$$E_6 > n_6$$

Replacing $E_6$ and $n_6$ in these conditions leads to the form

If $A_6^2 > 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6$.

If $A_6^2 < 2Y_6$ then $A_6 > (X_6 - Y_6)/2M_6 .+ (Y_6 - A_6^2/2)^2/2M_6(X_6 - Y_6)$.

Note that by normalizing $A_6$ by $M_6$ and $X_6$ and $Y_6$ by $W_6^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_6^2 + (M_6 + U_6)\lambda_6 - (M_6 + V_6)\lambda_6 e^{2T_6\lambda_6} + (M_6 U_6 + Y_6)$$

$$-[M_6(U_6 + V_6) - X_6]\, e^{2T_6\lambda_6} + M_6 V_6 e^{2T_6\lambda_6} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_6 = \mathrm{i} Z_6$ *and* separate

(A3.1)into real and imaginary parts

$$Z_6^2 + (M_6 + V_6) Z_6 \sin(Z_6 T_6) - (U_6 M_6 + Y_6) + (M_6(U_6 + V_6) - X_6)$$

$$\mathrm{X}\cos(Z_6 T_6) - 2M_6 V_6 \cos^2(Z_6 T_6) + M_6 V_6 = 0$$

$$(M_6 + U_6) Z_6 - (M_6 + V_6) Z_6 \cos(Z_6 T_6) + (M_6(U_6 + V_6) - X)$$

$$\mathrm{X}\sin((Z_6 T_6) - M_6 V_6 \sin(Z_6 T_6) = 0$$

Using (A.3.3) to eliminate $Z_6$ and $Z_6^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_6 =$

$1 - \cos Z_6 T_6$ and $A_6 = U_6 - V_6$ we obtain

$$4M_6^2 V_6^2 W_6^4 + 2M_6 V_6 W_6^3[(M_6 - V_6)^2 + A_6(M_6 + V_6) + 2(A_6 M_6 - X_6)]$$

$$+ W_6^2 [M_6 - V_6)^2 (2 A_6 M_6 + Y_6 - X_6) + (2A_6 M_6 - X_6)^2 - A_6 (M_6 - V_6)(2A_6 M_6 - X_6)$$

$$+ M_6 V_6 (A_6^2 + 4(X_6 + Y_6))]$$

$$+ W_6 [A_6^2 + 2X_6)(2 A_6 M_6 - X_6) + 2A_6 Y_6 (M_6 + V_6) - M_6 A_6^3] + A_6^2 (X_6 M_6 - Y_6) = 0$$

Being a quartic in $W_6$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_6 > 0$$

$$and$$

$$C_6 > (X_6 - Y_6)/2 M_6$$

$recalling\ that\ A_6 = U_6$ and defining $C_6 = U_6 + V_6$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_6 (0, U_6 V_6 = A_6^2 (X_6 + Y_6)$$

$$Q_6 (2, U_6 V_6) = (C_6 + 2M_6^2)^2 (2C_6 M_6 + Y_6 - X_6)$$

So that given (A3.6) the extreme $W_6$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_6$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_6$ the existence of a D-position at this point is

determined by letting cos $(Z_6 T_6) = -1\ in\ (A3.3)$

ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 e^{-\tau_6\lambda_6} + X_6 + Y_6 = 0 \qquad (A1.1)$$

Where $V_6 = dB_6/dR_6$ $U_6 = dD_6/dR_6 + N_6 dG_6/dR_6$ $X_6 = N_6G_6dF_6/dR_6$ and $Y_6 = -N_6G_6dM_6/dR_6$ (evaluated at the equilibrium). Note that $U_6X_6$ and $Y_6$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_6$ as $(T_6)$ . (It follows that $f(T_6)$ 0 for all $T$, and $\int_0^\infty f(T_6)dt = 1.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_6^2 + U_6\lambda_6 - V_6\lambda_6 \int_0^\infty e^{-\tau_6\lambda_6} f(T_6) + X_6 + Y_6 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MODULE NUMBERED TWO**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed

the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).. Haimovici's and Nunny's models have been applied to some of the Deleuzean and Heideggerian problems too : Scoring on that count are :

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time

delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t - T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if

only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

TIME DELAY

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**MATURATION AND CONSUMMATION OF CONFEDERATION OF DEBIT AND CREDIT TYPES:**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t-T_7)F_7(R_7(t-T_7) - N_7(t)M_7((R_7(t) \qquad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7} + \frac{N_7 dG_7}{dR_7} - \frac{dB_7}{dR_7} > \frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7} + \frac{dM_7}{dR_7}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes $(L_7)$ and the other derived from the non-delayed or "immediate" processes$(I_7)$:as in payment of usane bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_7 = \frac{N_7 G_7}{2M_7}\frac{d(7nF_7)}{dR_7} (\geq 0) \qquad (7a)$$

$$I_7 = -\frac{N_7 G_7}{2M_7}\frac{d(7nM_7}{dR_7} \ (\geq 0) \qquad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7 > L_7 - I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_7 > L_7 - I_7, \quad \text{when} \quad P_7^2/4 > I \qquad (8a)$$

$$P_7 > L_7 - I_7 + \frac{(I_7 - P_7^2}{(L_7 - I)} \quad \text{when} \quad P_7^2/4 < I. \qquad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

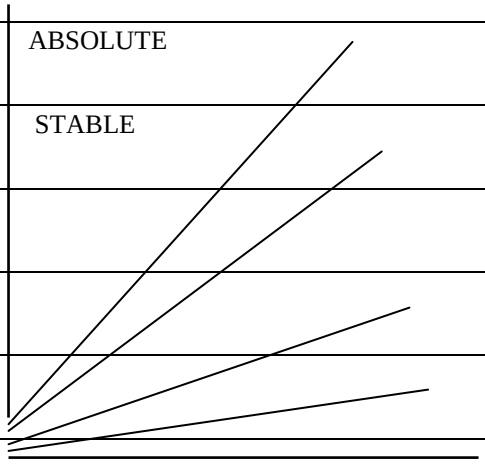



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit

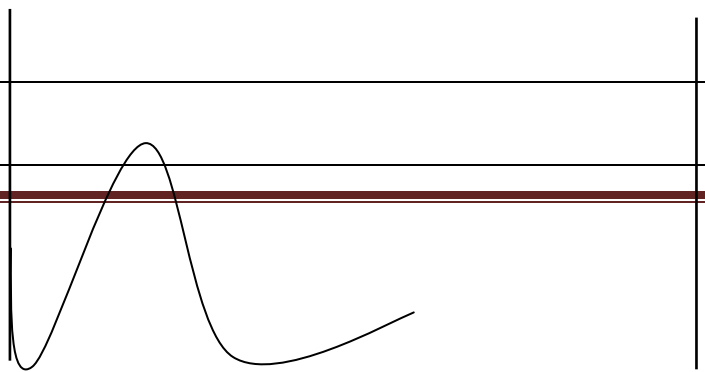
(like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

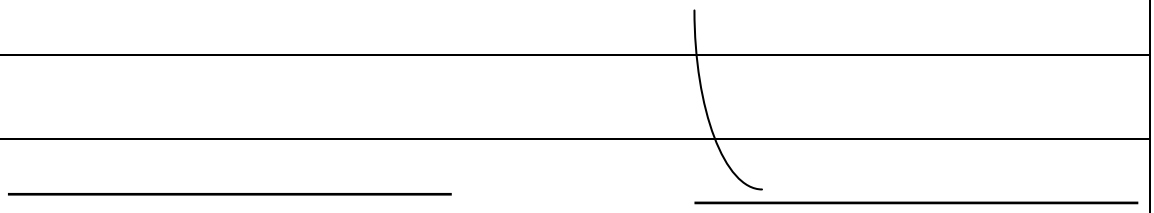
Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the





TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t-T_7)F_7\big(R_7(t-T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t-T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation $(dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7)$ and the effective prey regulation $(dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7)$ must be positive.

The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the

total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

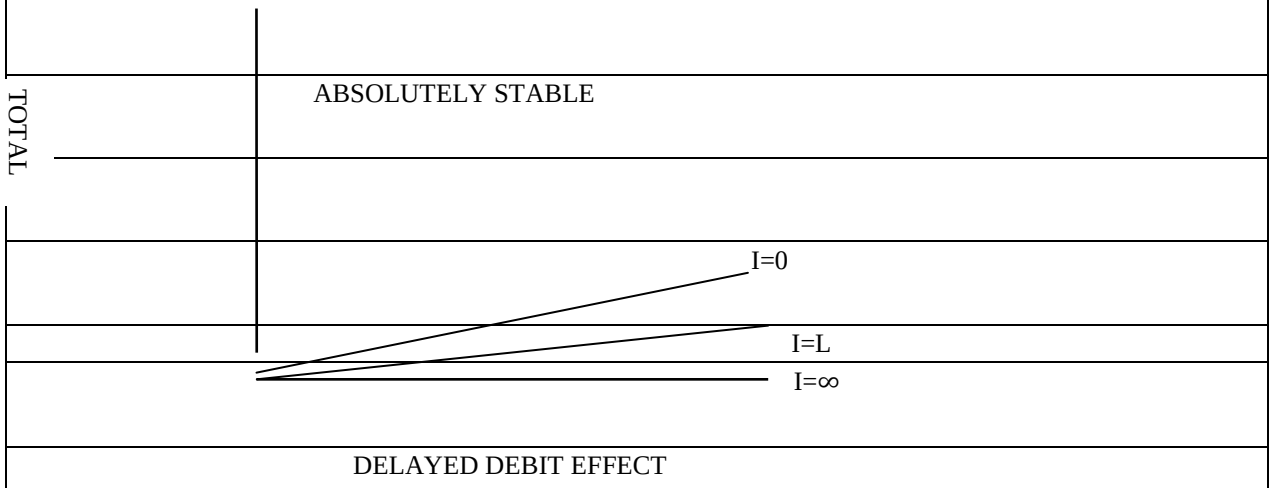


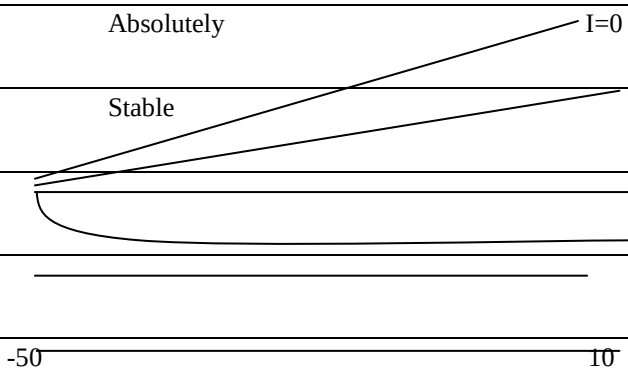

FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.






**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay.

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^\infty \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7\left(U_7 - V_7\int_0^\infty cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7\right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7\int_0^\infty cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$then given the constraints of $f(T_7)$, $(A1.5)$*can never be satisfied Co*nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7\lambda_7 e^{-T_7\lambda_7} + [(U_7 - V_7)M_7 + Y_7] - [((U_7 -$$

V 7M 7−]$e^{-T7\lambda 7}$=0

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$Z_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7)\cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7)\sin(Z_7 T_7) = 0$$

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $Z_7^2$ and $Z_7$in (A2.2):

$W_7^2 (X_7^2 - 3 M_7 A_7 X_7 - M_7^2 (X_7 - Y_7) + 2 M_7^3 A_7 + 2 A_7^2 M_7^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2(X + Y) = 0$$

*where* $W_7$=1-cos $(Z_7 T_7)$.

In the absence of a delay the system is stable for $A_7 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7 > (X_7 - Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$$A_7 > [\ X_7 - Y_7/2\ M_7$$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_7, A_7) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_7 \le 2$ *is* of interest First consider the extreme values

$$Q_7(0, A_7) = A_7^2\ (X_7 - Y_7)$$

$$Q_7(2, A_7) = (A_7 + 2\,M_7)^2\,(2A_7 M_7 - (X_7 - Y_7))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7 M_7^2 - 3M_7 X_7\right) + W_7\left(3M_7 A_7^2 - 2A_7 X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7+Y_7\right)$$

$=2M_7^2W_7^2(M_7+\mathrm{A}_7)+M_7A_7^2W_7+\left(2A_7M_7-X_7+Y_7\right)(A_7W_7+M_7W_7^2$

$+(2-W_7)((2M_7X_7+M_7Y_7)W_7+A_7(X_7+Y_7)),$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $W_7$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$has no minimum in the range $0 \le W_7 \le 2$.Defining

$Q_7\left(W_7,A_7\right)=a_7W_7^2+c_7 where\ a_7\ b_7\ and\ c_7$ Are functions of $A_7$we can define the condition for a minimum in the range 0-2

condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7 \ge b_7 \ge 0$$

And

$$b_7^2 \le 4a_7c_7$$

$$Case(i): Y_7=0$$

$Let\ Y_7=0$ And try the substitution A= $X_7=/2M_7\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}\ (W_7 X_7 / 2M_7)$ is positive for *W<2. Since* $\mathrm{Q}_7$ increases with $A_7$(see A2.9), increasing $A_7$will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

$Case\ (ii): Y_7 = y_7\ X_7$ ,

$Let\ Y_7 = y_7\ X_7$And try the substitution $A_7 = X_7/(2 + A_7\ y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + n_7 y_7)$ we have

$$W_7^2\ y_7\left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2\left(2 - n_7 + n_7 y_7\right)\right]$$

$$- W_7\ X_7\left[\frac{X_7(1 + n_7 y_7)}{M_7^2(2 + n_7 y_7)^2} + 2\ y_7\left(n_7 - 1\right) + \frac{X_7^2(7 + y_7)}{M_7^2(2 + n_7 y_7)}\right] = 0.\ (A2.12)$$

Since $n_7$ and $A_7$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7), the\ condition (A2.77?) can\ be\ written\ as$

$\left[\frac{X_7(E_7 - 7)}{E_7^2}\right]^2\left[\frac{X_7^2}{M_7^2} - 4\ y_7\ E_7^2\frac{X_7}{M_7^2} + 4y_7(n_7 - 2)E_7^3\right]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$2\ E_7[y_7 E_7 + (y_7\ E_7(7 + y_7)(E_7 - n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - y_7 E_7(7 + y_7)(E_7 - n_7))^{7/2}]$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2$) is large or small then (A2.6) will have a real solution although if may not be in the range (0 W 7 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7 - 7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2 - 2y_7\frac{(E_7 + n_7^2 y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right) - 4y_7(E_7 - n_7)$$ 0

And has the solution

$\frac{X_7}{M_7^2} > \frac{E_7}{(E_7-7)}[y_7(E_7 + n_7^2 y_7) + [y_7^2 (E_7 + n_7^2 y_7)^2 + 4 y_7 (E_7 - 7)(E_7 - n_7)]^{7/2}]$

Given $E_7 > n_7$ this lower limit to $X_7 / M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7 / M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$ in the appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7 < n_7$ $(ie., 2A_7 M_7^2 < X_7 - Y_7)$ then (A2.7)and (A2.8)show that $Q_7$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2} > 2E_7[y_7 E_7 - (y_7 E_7 (7 + y_7)(E_7 - n_7)]^{7/2}]$$

And

$$E_7 > n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2 > 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.

If $A_7^2 < 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7 + (Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$ and $X_7$ and $Y_7$ by $W_7^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSILIDATAN**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7\lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]e^{2T_7\lambda_7} + M_7 V_7 e^{2T_7\lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i}\, Z_7$ *and* separate

(A3.1)into real and imaginary parts

$$Z_7^2 + (M_7 + V_7)Z_7 \sin(Z_7 T_7) - (U_7 M_7 + Y_7) + (M_7(U_7 + V_7) - X_7)$$

$$\mathrm{X}\cos(Z_7 T_7) - 2M_7 V_7 \cos^2(Z_7 T_7) + M_7 V_7 = 0$$

$$(M_7 + U_7)Z_7 - (M_7 + V_7)Z_7 \cos(Z_7 T_7) + (M_7(U_7 + V_7) - X)$$

$$X \sin\left(\left(Z_7 T_7\right) - M_7 V_7 \sin\left(Z_7 T_7\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_7$ and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - cos\ Z_7\ T_7$ and $A_7 = U_7 - V_7$ we obtain

$$4 M_7^2 V_7^2 W_7^4 + 2 M_7 V_7\ W_7^3 [(M_7 - V_7)^2 + A_7 (M_7 + V_7) + 2(A_7 M_7 - X_7)]$$

$$+ W_7^2 [M_7 - V_7)^2 (2 A_7\ M_7 + Y_7 - X_7) + (2A_7\ M_7 - X_7)^2 - A_7 (M_7 - V_7)(2A_7\ M_7 - X_7)$$

$$+ M_7 V_7 (A_7^2 + 4(X_7 + Y_7))]$$

$$+ W_7 [A_7^2 + 2X_7)(2 A_7\ M_7 - X_7) + 2A_7 Y_7 (M_7 + V_7) - M_7 A_7^3] + A_7^2 (X_7\ M_7 - Y_7) = 0$$

Being a quartic in $W_7$ (A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7 > 0$$

$$and$$

$$C_7 > (X_7 - Y_7)/2 M_7$$

$recalling\ that\ A_7 = U_7$ and defining $C_7 = U_7 + V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0, U_7 V_7 = A_7^2 (X_7 + Y_7)$$

$$Q_7\left(2, U_7 V_7\right) = \left(C_7 + 2M_7^2\right)^2 (2C_7 M_7 + Y_7 - X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1$ *in* $(A3.3)$

## ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

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# Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

**Black holes Teviet Creighton and Richard H. Price (2008), Scholarpedia, 3(1):4277. doi:10.4249/scholarpedia.4277 revision #135523 [link to/cite this article] Post-publication activity Curator: Teviet Creighton Dr. Teviet Creighton, University of Texas, Brownsville, TX Richard H. Price, University of Texas, Brownsville, TX Figure 1: Artist's conception of an accretion disk around a black hole. Credit: NASA/CXC/M.Weiss**

(2717) In this way, stellar mass holes produce (eb) detectable X-rays, supermassive black holes produce *(e&eb) a wide spectrum of electromagnetic signals, and both types can be inferred from (e) the orbital motion of luminous stars and (e&eb) matter around them.

(2718) Phenomena involving black holes of any mass can produce (eb) strong gravitational waves, and are of interest as (=) sources for present and future gravitational wave detectors.

**Classical vs. relativistic black holes**

(2719) Something like a black hole exists within (eb) Newton's classical theory of gravity.

(2720) In that theory, an energy argument tells us that there is (=) an escape velocity vesc=2GM/R−−−−−−−√ from (e) the surface of any spherical object of (e) mass M and radius R

(2721) If this velocity is greater than (e) the speed of light c then (eb) light from this object cannot escape to infinity.

(2722) Thus the condition for such an "unseeable" object is R<2GM/c2

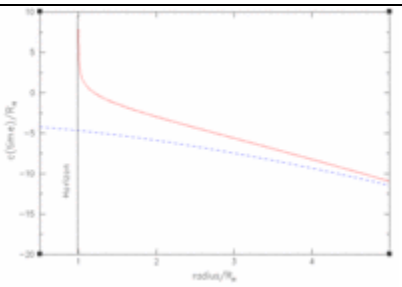
(2723) In the classical theory, a particle could overcome (e) this gravity with strong enough engines to provide (eb) the energy needed for escape.

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| (2724) This is not so in general relativity, Einstein's theory of gravitation. In that theory, escaping the black hole is (=) equivalent to moving faster than light, impossibility in relativity. |
| (2725) To understand the relativistic black hole it is useful to think of (e) space being dragged inward towards (e&eb) a gravitational center, at a faster rate near (eb) the center than far from it. |
| (2726) The distance at which space is moving inward at the speed of light represents (eb) the location of the event horizon, since (e) no signal can progress outward through (e&eb) space faster than c . |
| (2727) This comparison is more than a metaphor; black hole analog experiments with (e&eb) accelerating gas flows and (e&eb) other phenomena are being designed. |
| (2728) An important difference from Newton's theory is that Einstein's, and other relativistic theories of gravitation are (=) nonlinear in the sense that gravitation (as well as mass) can be (=) a source of gravity. |
| (2729) Thus when a massive object collapses (eb) small enough, the tendency to continue the collapse and form (eb) a black hole can become unstoppable. |

**Stationary black holes**

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| (2730) In the Newtonian theory, gravity is described by (e) the potential $\Phi$. |
| (2731) Inside a spherical object the form of (e) $\Phi(r)$ depends on (=) the interior structure, but in the vacuum outside (e) matter the potential $-GM/r$ depends only on (=) the interior mass. |
| (2732) Similarly, in Einstein's theory the stationary (time-independent) spherically symmetric exterior solution, called the Schwarzschild spacetime, depends only on (=) the mass of the interior object. |
| (2733) If the interior object is small enough, then the Schwarzschild exterior extends to (e&eb) small enough radius that there is (=) a horizon, a surface across which light cannot (e) move outward. |
| (2734) This horizon radius $RH=2GM/c2$ is, coincidentally, (=) the same as the critical radius for (e) "unseeable" objects in Newton's theory. |
| (2735) The meaning of "radius" as distance to (e) the center is not straightforward for (e) the Schwarzschild solution. |
| (2736) Radius RHhere actually means that the area of the event horizon is (=) $4\pi R2H$. |
| (2737) In Einstein's theory, the "exterior" solution can be taken to apply with no (e) interior solution. |
| (2738) In this case it is gravity itself, rather than matter, that acts as (=) the source of gravity. |
| (2739) The inward-extended exterior solution does not (e) reach a center, but rather is connected via (e&eb) a spacetime bridge to (e&eb) another universe, or another section of our own. |
| (2740) For an astrophysical black hole, formed from (e) the collapse of matter, a physical solution for (e) the matter distribution replaces (e&eb) the pure vacuum Schwarzschild solution in (eb) the interior of the black hole. |
| (2741) This physical solution lacks (e) the Spacetime Bridge of ideal mathematical black holes, but contains (e) a central "singularity" where (e) matter is compressed to infinite density. |
| (2742) Very close to this singularity it is expected that the laws of general relativity will no longer apply, and as-yet unknown laws of quantum gravity are needed. |
| (2743) A more general stationary black hole solution of (e) Einstein's theory is the Kerr solution, a vacuum spacetime with (e&eb) both mass and angular momentum, and taken to represent (eb) a rotating black hole. |

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| (2744) In its pure mathematical form the Kerr hole contains (e) a spacetime bridge, but as in the case of the Schwarzschild black hole this bridge is absent in (eb) realistic black holes that form by (e) the collapse of matter. |
| (2745) Unlike the Schwarzschild spacetime, the Kerr solution is not the exterior spacetime of (e) a material object with (e&eb) angular momentum. |
| (2746) In fact no realistic solution has been found to join a Kerr exterior to (e&eb) a material interior. |
| (2747) The Kerr solution only becomes (=) the exterior spacetime asymptotically at (eb) very late times after (e) the collapse of an object. |
| (2748) Two other exact mathematical black hole solutions are (=) the Reissner-Nordström spacetime, representing (eb) a hole with mass and electrical charge, and the Kerr-Newman spacetime, representing (eb) a hole with mass, electrical charge, and angular momentum. |
| (2749) These spacetimes are not astrophysically relevant, since (e) astrophysical bodies have (e) negligible net electrical charge. [For a detailed description of these spacetimes see, e.g., Misner et al. (1973), Part VII; or Wald (1984), Chap. 12.] |
| (2750) All of these spacetimes, including (e) the ones with angular momentum, are stationary: that is, they are (=) independent of time. |
| (2751) But in relativity there is (=) no unique meaning to time, so an important question is: "Just what 'time' is it of (e) which the stationary black holes are independent?" |
| (2752) The answer lies in the fact that one can assign (e&eb) every spacetime point four coordinates, four labels that uniquely identify (e&eb) the location of each point. |
| (2753) One of these coordinates is (=) called the "coordinate time." |
| (2754) Spacetimes that are said to be stationary, like the spacetime of a Kerr hole, have (e) a special property: the time coordinate may be chosen so that (e) the spacetime geometry is (=) the same at any moment of this time coordinate. |



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| (2755) Figure 2: Coordinate time vs. (e&eb) proper time for the radial fall of (e) a particle towards a Schwarzschild black hole |
| (2756) At large distances from a stationary black hole, where (e) spacetime curvatures are weak, this stationary time coordinate can be chosen also to have (e) another important property: to agree with (=) the "proper time," or ordinary clock time, of (e) an observer at rest with respect to (e&eb) the hole. |
| (2757) Since we ourselves are more-or-less at rest (or are at (eb) nonrelativistic velocities) very far from black holes, this kind of stationary coordinate time is (=) the time used in astronomical observations. |
| (2758) For observers near the hole, however, proper time and (e&eb) stationary coordinate time will not agree |

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| (2759) An interval of proper time between two events is (=) shorter (near the horizon, much shorter) than the interval of stationary coordinate time between (e&eb) those two events. |
| (2760) The relationship between stationary coordinate and (e&eb) proper time is further complicated by (e) special relativistic time dilation for (e&eb) observers who are rapidly moving. |
| (2761) Figure 2 illustrates this by showing the coordinate time vs. radius for a particle falling into (eb) a black hole, and comparing it with (e&eb) the proper time measured by (e) an observer riding along with (e&eb) the infalling particle. |
| (2762) The progress as measured with proper time is in no way special at (eb) the horizon: falling observers will not notice anything strange as they pass (e&eb) the point of no return. |
| (2763) But as described in coordinate time, the observer (and likewise, the surface of a collapsing star) takes an infinite amount of time to reach (eb) the horizon. |
| (2764) Since coordinate time is the proper time of (e) distant observers, astronomers will see the particle reach (eb) the horizon only in the infinite future. |
| (2765) The two (or more) types of time are sometimes a source of (e) confusion in discussions of black hole phenomena, since they often give totally different answers to the question "how long does it take?" |

**Black hole parameters**

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| (2766) Astrophysical black holes are characterized by (e) two parameters: their mass and their angular momentum (or spin). |
| (2767) The mass parameter M is equivalent to (=) a characteristic length GM/c2=1.48km (M/Mo) , or a characteristic timescale GM/c3=4.93×10−6sec(M/Mo) , where (e) Mo denotes the mass of the Sun. |
| (2768) These scales, for example, give (eb) the order of magnitude of the radii and (e&eb) periods of near-hole orbits. |
| (2769) The timescale also applies to (e&eb) the process in which a developing horizon settles into (e&eb) its asymptotically stationary form. |
| (2770) For a stellar mass hole this is of order 10−5sec, while for (e) a supermassive hole of 108Mo, it is thousands of seconds. |
| (2771) For Schwarzschild holes, and approximately for Kerr holes, the horizon is (=) at radius RH=2GM/c2. |
| (2772) At the horizon, the "acceleration of gravity" has no meaning, since (e) a falling observer cannot stop at the horizon to be weighed. |
| (2773) What are relevant at the horizon are (e) the tidal stresses that stretch and distort (e&eb) the falling observer. |
| (2774) This tidal stretching is given by (e) the same expression, the gradient of the gravitational acceleration, as in Newtonian theory 2GM/R3H=c6/(4G2M2) |
| (2775) In the case of a solar mass black hole the tidal stress (acceleration per unit length) is (=) enormous at the horizon, on the order of 3×109(Mo/M)2s−2 : that is, a person would experience (e) a differential gravitational field of about 109 Earth gravities, enough to rip apart (e&eb) ordinary materials. |
| (2776) For a supermassive hole, by contrast, the tidal force at the horizon is smaller by (e) a typical factor 1010--16 and would be easily survivable |
| (2777) However, at the central singularity, deep inside the event horizon, the tidal stress is (=) infinite. |
| (2778) In addition to its mass M, the Kerr spacetime is described with (e&eb) a spin parameter a defined by (e) the dimensionless expression aM=cJGM2, where J is the angular momentum of the hole. |

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| (2779) For the Sun (based on (e) surface rotation) this number is about 0.2, and is much larger for many stars. |
| (2780) Since angular momentum is ubiquitous in astrophysics, and since it is expected to be (=) approximately conserved during (e&eb) collapse and black hole formation, astrophysical holes are expected to have (e) significant values of a/M, from (e) several tenths up to and approaching unity. |
| (2781) The value of a/M can be (=) unity (an "extreme" Kerr hole), but it cannot be (=) greater than unity. |
| (2782) In the mathematics of general relativity, exceeding this limit replaces the event horizon with (e&eb) an inner boundary on the spacetime where tidal forces become (=) infinite. |
| (2783) Because this singularity is (=) "visible" to observers, rather than (eb) hidden behind a horizon, as in a black hole, it is called a naked singularity. |
| (2784) Toy models and heuristic arguments suggest (eb) that as a/M approaches unity it becomes (=) more and more difficult to add angular momentum. |
| (2785) The conjecture that such mechanisms will always keep a/M below (eb) unity is called cosmic censorship. |
| (2786) The inclusion of angular momentum changes details of the description of the horizon, so that, for example, the horizon area becomes (=) Horizon area= $(4\pi G2/c4)[(M+M2-a2\sqrt{\phantom{xx}})2+a2]$. |
| (2787) This modification of the Schwarzschild (a=0) result is not significant until (e) a/M becomes very close to (e) unity. For this reason, good estimates can be made in many astrophysical scenarios with a ignored. |

**Dynamical black holes**

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| (2788) The event horizon is (=) defined as the outer boundary of the region from which there is (=) no escape. |
| (2789) For stationary black holes this surface is at (eb) a fixed location in space, but more generally the horizon is (=) dynamical; it can grow, change shape, and oscillate. In particular, **a horizon can be born**. |
| (2790) The birth of a horizon is a change first studied by Oppenheimer and Snyder for the collapse of spherical pressure less fluid. |
| (2791) The general scenario is for a horizon to be (=) born deep inside the collapsing matter and to spread outward. |
| (2792) Small changes in a horizon can be treated as (=) perturbations, greatly simplifying the mathematics of Einstein's equations. |
| (2793) In this way, it has been found that black holes have (e) characteristic patterns of oscillations, called quasinormal modes. |
| (2794) These modes are (=) like mechanical resonances, but are (=) highly damped by the emission of (e) gravitational waves, and have (e) both periods and damping times on the order of (e) the characteristic time GM/c3. |
| (2795) Figure 3: Distorted horizons of a black hole binary near (eb) the point of merger: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/ |
| (2796) For large nonspherical changes in (eb) a dynamical horizon, only supercomputer simulations can give (eb) quantitative answers. |
| (2797) The focus of such work has been the merger of (e) two black holes in binary orbits around (e&eb) each other, a scenario of special interest as (=) a source of strong gravitational wave emission. |

(2798) Only in 2005 were technical problems first overcome [Pretorius (2005) (See Fig. 3.), Campanelli et al. (2006), Baker et al. (2006)] so that an accurate picture could emerge of how two horizons join to become (=) a single final horizon.

(2799) There is one change that a horizon cannot make according to classical (nonquantum) general relativity: it cannot (e) decrease its area.

(2800) But considerations by Stephen Hawking and others of quantum effects in black hole spacetimes suggest that radiation arising in (eb) the close exterior of the black hole can carry off (e&eb) energy, and decreases (e) the mass (and hence horizon area) of the hole.

(2801) Although no quantum theory of relativistic gravitation currently exists, it is generally accepted that this Hawking radiation will be (=) a feature of any such theory.

(2802) The radiation behaves as if the horizon were (=) a blackbody (perfect thermal emitter) at a temperature $6\times10-8$ (Mo/M) K.

(2803) Thus for astrophysical black holes, mass loss by (e) Hawking radiation is much less important than the mass increase due to (e) absorption of the 3K cosmic microwave background, and that mass increase is (=) itself negligible.

**Astrophysical black holes**

(2804) Black holes in our Universe can be grouped as: primordial black holes, stellar-mass black holes, and supermassive black holes. The first is highly speculative; the second and third are broadly accepted.

(2805) Primordial black holes of all masses are postulated to have formed from (e) quantum fluctuations in the early Universe, but those with (e&eb) mass less than around 1013kg would have already evaporated due to Hawking radiation.

(2806) No significant observational evidence has yet been found of (e) the existence of these objects [MacGibbon and Carr (1991)].

(2807) Stellar-mass black holes, ranging from a few to a few tens of solar masses, are (=) normal but rare endpoints in (eb) the evolution of massive stars.

(2808) When a star exhausts its nuclear fuel and cools, it must collapse unless (e) it is supported by nonthermal forces.

(2809) It is known that such nonthermal forces cannot (e) resist gravitational compression for (e) masses greater than around 1.5 to 3Mo. (The uncertainty is due to (e) uncertainty in our knowledge of nuclear physics at high densities.)

(2810) There are many stars far more massive than this, but in their death throes massive stars expel (eb) a great deal of their mass in supernova explosions

(2811) Even stars initially as massive as 20-30Mo may blow off enough mass to (e) leave a remnant neutron star smaller than 1.5Mo.

(2812) Stars more massive than 20-30Mo may form (eb) a neutron star core, only to have it collapse due to (e) fallback of material from (e) the stellar mantle. Still more massive stars (above ~40Mo) may form (eb) a black hole directly in collapse, with or without (e&eb) a supernova explosion.

(2813) These masses are (=) rather uncertain at present, as our understanding of supernova explosion mechanisms is (=) still evolving [Woosley and Janka (2005), Muno (2006)].

(2814) Collapsing gas clouds larger than ~100Mo are expected to dissipate from (e) radiation pressure, which would prevent (e) more massive stars from forming, and consequently set (eb) an approximate upper limit for stellar mass black holes.

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| (2815) Observational evidence for stellar-mass black holes comes primarily from (e) X-ray astronomy. |
| (2816) A black hole in a close binary orbit can form (eb) an accretion disk (see Figure 1) of matter pulled off (er) a normal stellar companion. In its inspiral, disk material is heated by shearing, and becomes (=) a strong X-ray emitter. |
| (2817) If observations reveal a point-like X-ray emitter with (e&eb) a mass inferred from (e) orbital dynamics to be (=) above that possible for a neutron star, it becomes (=) a black hole candidate. |
| (2818) t present there are over a dozen such black hole candidates, including the first object to be identified as (=) a black hole, the X-ray source Cygnus X-1 (estimated mass 10±3Mo) [Casares (2006)]. |
| (2819) The other class of observationally-supported black holes is (=) that of supermassive black holes, ranging from hundreds to (e&eb) billions of solar masses. |
| (2820) Evidence for the existence of black holes at (eb) the upper end of this range is overwhelming. |
| (2821) By contrast, the existence of black holes with (e&*eb) mass roughly of order 102 to 104Mo, referred to as "intermediate-mass black holes," is speculative. |
| (2822) The supermassive holes may have begun as (=) primordial or stellar-mass black holes, but have grown through (e&eb) absorption of stars or gas or through mergers with (e&eb) other holes [Ferrarese and Ford (2005)]. |
| Figure 4: Observed orbits of stars near (eb) the dark compact object at the center of (e) our Milky Way galaxy: still frame from an animation. Animation created by Prof. Andrea Ghez and her research team at UCLA, from data sets obtained with the W. M. Keck Telescopes: http://www.astro.ucla.edu/~ghezgroup/gc/ |
| Evidence for supermassive black holes originally came from (e) quasars, small but intense radio sources seen at (eb) cosmological distances. |
| Their huge luminosities implied (eb) high mass, while (e) their rapid variability implied (eb) extremely small size. |
| More recently, measurements of Doppler shifts of stars and gas in the centers of galaxies have shown (eb) that compact objects of mass greater than 106Mo reside in (eb) the cores of most galaxies; the very small size of these central objects rules out any plausible alternative to the black hole explanation. |
| The black hole at the center of our own Milky Way galaxy has (e) a measured mass of ∼3×106Mo , based on (e) the orbits of stars near the radio source Sagittarius A* associated with (e&eb) the Galactic core [Ghez et al. (2005); see Fig. 4]. |
| (2828) Black holes are of particular interest for (e) gravitational waves, and vice versa. |
| (2829) For gravitationally-bound systems, the typical maximum gravitational wave amplitude (dimensionless strain) is (=) h∼(2GM/c2)2/(RD)∼R2H/(RD) , where D is the distance to the system, M and R are the mass and size of (e) the system, and RH is the horizon radius of (e) a black hole of the same mass. |
| (2830) From this it is clear that strongest gravitational waves will involve (e&eb) systems at or near black hole compactness: in particular, systems containing (er) black holes in close proximity to one another [Flanagan and Hughes (2005)]. |
| (2831) Conversely, a system containing only black holes would not be expected to radiate (eb) any sort of radiation other than gravitational waves, and in any system containing (e) black holes, gravitational waves provide (eb) the most accurate probes of conditions deep in (eb) their gravitational wells. |
| (2832) In particular, gravitational wave observations could definitively test (e&eb) whether a compact mass is truly a black hole [Cutler and Thorne (2002)]. |

(2833) The strongest intrinsic source of gravitational waves in the present Universe is (=) the collision of two comparable-massed black holes.

(2834) A sufficiently tight black hole binary system will lose (e) energy through emission of (e&eb) gravitational waves, causing (eb) the orbit to circularize and then slowly shrink over time.

(2835) Figure 5: Gravitational waves from (e) a black hole binary inspiral: still frame from (e) an inspiral movie. Numerical simulation and animation produced by Frans Pretorius, http://physics.princeton.edu/~fpretori/

(2836) When the orbital radius approaches (eb) the Schwarzschild radius of the system, complex nonlinear dynamics come into play (e&eb) and the final stages of (e) the inspiral and merger can only be modeled through (e&eb) numerical simulations.

(2837) These simulations are quite challenging, but recent breakthroughs [Pretorius (2005)(See Fig. 5), Campanelli et al.(2006), Baker et al.(2006)] have led to (eb) the first complete inspiral-merger gravitational waveforms.

(2838) After the black holes have coalesced, they form (eb) a single highly-distorted black hole which quickly settles down to (e) a quiescent Kerr state through (e&eb) mission of quasinormal "ringdown" gravitational waves [Pretorius (2005), Campanelli et al. (2006), Baker et al. (2006)].

(2839) Another gravitational emission mechanism for supermassive black holes would come from (e) their capture of compact stellar objects, particularly stellar-mass black holes.

(2840) Unlike comparable-mass mergers, these captures would be on (eb) highly eccentric orbits that would not circularize before merger, resulting in (eb) intricate rapidly precessing orbits in (eb) the final months or years of inspiral.

**(2841)** The consequent gravitational-wave signal would allow (eb) for exquisitely precise measurements of the parameters of (e) the supermassive black hole, as well as tests of any deviation from (e) the predictions of general relativity [Cutler and Thorne (2002)].

**(2842)** For stellar-mass black holes (tens of Mo), the final stages of inspiral and merger occur at (eb) frequencies of hundreds to thousands of Hz, a regime targeted by (e) ground-based gravitational-wave observatories.

**(2843)** Systems of supermassive black holes that are (=) millions of Mo emit at frequencies of tens of (e) mHz, which will be detectable by (e) proposed space-based (e) gravitational-wave detectors.

**(2844)** Black holes in the billions of Mo, or a stochastic background of signals from (e) more moderate supermassive black holes at (eb) earlier stages of their inspiral, might emit (eb) waves at nHz frequencies that would be measurable through (e&eb) years-long correlated timing measurements of radio pulsars [Jenet et al. (2005)].

**An equation to quantify the origins of life on other planets July 5, 2016 by Bob Yirka report**

(2845) (Phys.org)—A pair of researchers, one with the Columbia Astrobiology Center in New York, the other with the University of Glasgow in the U.K. has come up with a mathematical equation that when (e) solved is meant to offer (eb) a means for estimating how often life begins on other planets. In their paper published in Proceedings of the National Academy of Sciences, Caleb Scharf and Leroy Cronin describe their equation, how they came up with it and why they believe it might become more useful as scientists learn more about the true nature of other planets and solar systems.

(2846) Back in the 1960's, an astronomer named Frank Drake came up with a formula for (e) estimating the likely number of (e) alien civilizations that might be capable of transmitting radio signals in such a way as to be (=) recognizable by receivers here on earth.

(2847) That equation, named quite naturally the **Drake Equation** has been the unofficial standard-bearer for half a century, despite it having no solution because (e) some of the parameters are still unknown.

(2848) The same holds true for the new equation developed by (e) Scharf and Cronin, but the results given should the unknowns ever be discovered would offer a different sort of answer than Drake was looking for.

(2849) To build their formula, the duo came up with several parameters to solve for (e) **abiogenesis**, which is the likelihood of an event occurring that leads to (eb) life beginning, they included (e) the number of possible building blocks, the mean number of such blocks per possible organism, the availability of building blocks that might exist during a given time period, expressed as a fraction and the probability that the existence of the building blocks would actually lead to (eb) life starting in a given unit of time. It looks like this: **Nabiogenesis (t)= Nb · 1/no · fc · Pa · t**

(2850) In essence the formula is meant to suggest (eb) that the probability of life beginning on a given planet is very likely connected to (e&eb) whether there are building blocks available on a given planet, and how much of them there might be.

(2851) As more is learned by space researchers it is hoped that the formula could help narrow down (e) search targets by offering (eb) a statistical probability of success for (e) a given planet.

(2852) The researchers note that the building blocks available do not necessarily have to be (=) of the type that led to (eb) life beginning on our planet.

**(2853)** They also suggest that implied in the formula is (=) the likelihood that life beginning events are (=) more likely to occur in solar systems where (e) there are multiple planets, which allow (eb) for an opportunity to share materials that could lead to (eb)building blocks **More information: Quantifying the origins of life on a planetary scale, Proceedings of the National Academy of Sciences, www.pnas.org/cgi/doi/10.1073/pnas.1523233113**

**Abstract**

(2854) A simple, heuristic formula with (e&eb) parallels to the Drake Equation is introduced to help focus discussion on open questions for (e) the origins of life in a planetary context.

(2855) This approach indicates (eb) a number of areas where (e) quantitative progress can be made on parameter estimation for (e) determining origins of life probabilities, based on (e) constraints from Bayesian approaches.

(2856) Authors discuss a variety of "microscale" factors and their role in determining (eb) "macroscale" abiogenesis probabilities on (e&eb) suitable planets.

(2857) Authors also propose that impact ejecta exchange between (e&eb) planets with parallel chemistries and (e&eb) chemical evolution could in principle amplify (eb) the development of (e) molecular complexity and (e&eb) abiogenesis probabilities.

(2858) This amplification could be very significant, and both bias conclusions about abiogenesis probabilities based on (e) the Earth and provide (eb) a major source of variance in the probability of life arising in (eb) planetary systems.

**(2859)** Authors use heuristic formula to suggest (eb) a number of observational routes for improving constraints on (e) origins of life probabilities. **Read more at: http://phys.org/news/2016-07-equation-quantify-life-planets.html#jCp**

**Brian Koberlein Shared publicly - Yesterday 8:00 PM the Illusion of Truth**

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| (2860) | Physical theories are often presented as (=) a description of what's really going on. |
| (2861) | Forces act on (e&eb) a baseball, causing (eb) baseball to fall. |
| (2862) | Atoms collide and fuse in (eb) the heart of a star, releasing heat and energy. Science is true, as is often said, and our scientific theories encapsulate this truth. But this isn't entirely true. |
| (2863) | For example, quantum mechanics is a strange, sometimes confusing theory. Objects can be (=) particles and waves. |
| (2864) | They do strange things like tunnel through (e&eb) barriers, and can even appear and disappear (e&eb). |
| (2865) | Given all this strangeness, what's really going on? What's the true nature of quantum reality? It depends on (=) which approach you want to take. |
| (2866) | In one view quantum objects are described by (e) a wave-like probability function. |
| (2867) | When you interact with (e&eb) or "observe" (e) the object this "wavefunction" collapses into (eb) a definite state. |
| (2868) | That's view commonly presented, but a quantum system can also be described by (e) its transitions between energy states. |
| (2869) | Since this uses a mathematical method involving (e&eb) matrices, it's known as (=) matrix mechanics. |
| (2870) | Both wavefunctions and matrix mechanics give (eb) the same results, but their view of what really goes on is very different. |
| (2871) | Then there's the path integral method |
| (2872) | . Rather than a wavefunction or matrix transitions, path integrals imagine quantum objects can take almost any path between (e&eb) two states. |
| (2873) | By summing all the possible paths you can derive (eb) the odds that it will occur. |
| (2874) | So which one is true? Are quantum objects distributed waves of (e) probability? |
| (2875) | Are they simply transitions between (e&eb) energy states? |
| (2876) | Do they take an infinite number of (e) paths between states? |
| (2877) | Each of these models makes (eb) the same predictions, so one version is no more "true" than the others. |
| (2878) | What happens in practice is that we'll use whatever method is useful at the time. They are equivalent models, so the best model for the job is the one we'll use. The only reason the wavefunction view is so common is that it's (=) the version usually taught to introductory students. You might think this uncertainty of truth is due to (e) the behavior of quantum physics itself. It's so strange and counterintuitive that we can't (e) wrap our puny brains around what's really going on. But the same thing occurs in lots of other fields. Even something as straight forward as basic Newtonian physics |
| (2879) | Toss a baseball in the air and (eb) Earth's gravitational force pulls it down. |
| (2880) | The force of gravity is a simple truth, right? While we often describe classical motion in terms of (e&eb) forces and acceleration, we can also describe it in terms of (e&eb) energy and momentum. |
| (2881) | In the Lagrangian and Hamiltonian approach, the path of a baseball is (=) the optimized path among possibilities. |
| (2882) | In this view a baseball's path is (=) the extrema of an energy equation, and force can be derived as a necessary consequence of (e) this. |
| (2883) | In the relativistic view the baseball follows (e&eb) a geodesic, which is (=) the minimal path through (e&eb) space and time. |

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| (2884) So is a baseball's motion due to (e) a gravitational force, an energy extrema, or a spacetime geodesic? |
| (2885) As with quantum theory, different approaches yield (eb) the same result. They are mathematically equivalent, so we can use whatever method is most useful at the time. |
| (2886) At its core, science is less about truth and more about models. The metaphysics underlying a model is useful only as far as it allows us (eb) to make better predictions, generate (eb) new ideas, or bring models together as (=) a cohesive whole |
| (2887) This is why we have no problem using (e) classical gravity to calculate the path of (e) a spacecraft through (e&eb) the solar system, while (e) using special relativity to account for (e) the Doppler shift of the spacecraft's radio signals. |
| (2888) It's why we can use (e) quantum physics to study atoms in the morning, and general relativity to (e) study black holes in the afternoon. |
| (2889) In regimes where models conflict with each other it isn't a failure of truth, but instead shows an opportunity to develop a better model. |
| (2890) It could be that with each better model we move closer to (e&eb) the truth about reality. |
| (2891) The truth is out there, and science strives to move towards (e&eb) that truth. |
| (2892) It's a common view, and certainly the search for truth has driven (eb) many scientists to develop better and better models. |
| (2893) But the real power of science is the recognition that what we have are models. Our models can be powerful, but they are always a bit tentative. There's always a chance that they might just be an illusion of truth. |

**Synopsis: A Single-Photon Cheshire Cat July 1, 2016 Researchers detected the polarization of a photon separate from the photon itself, just as the grin of Lewis Carroll's Cheshire cat can appear apart from the cat's body Synopsis figure APS/Carin Cain**

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| (2894) In a quantum Cheshire cat, one property of a particle is present detached from (e) the particle itself, just as (=) the grin of the feline in Alice in Wonderland can exist in the absence of (e) the whole animal. |
| (2895) Now a research team has successfully separated (e&eb) a single photon from (e) its polarization. |
| (2896) The effect was previously demonstrated (eb) with neutrons and also with light in a classical scenario, but the latest experiment uses single photons. |
| (2897) To produce the Cheshire cat, **Maximilian Schlosshauer and his colleagues** at the University of Portland in Oregon used an interferometer where (e) light hits a beam splitter, goes in two, opposite directions around a circuit (but along slightly offset routes), and then recombines at (eb) the entrance. |
| (2898) The incoming polarization of light was horizontal, but this was switched to (e&eb) vertical for the counterclockwise loop. |
| (2899) Each path included a device to adjust (e&eb) the polarization further and another to (e&eb) scatter away a small number of photons of (e) a designated polarization, the equivalent of detecting their presence. |
| (2900) The team fired single photons into (e&eb) the interferometer so that only one photon was present at (eb) a time, traversing both paths simultaneously. |
| (2901) They placed a horizontal polarizer in front of the output photon detector and gradually varied (e&eb) the difference between the two path lengths. |

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| (2902) This variation generated (eb) a sine-wave interference pattern from (e) the exiting photons only when the polarization was tweaked in (eb) the counterclockwise path. |
| (2903) At the same time, the average number of photons counted decreased only when (e) light was scattered from the clockwise path. |
| (2904) Simultaneously demonstrating these different effects for two routes of a single photon constitutes (eb) evidence of the Cheshire cat. |
| **(2905)** The effect might someday improve precision experiments if (e) an unwanted property could be moved away from the property being measured. **This research is published in Physical Review A. – David Ehrenstein David Ehrenstein is the Focus Editor for Physics.** |

**Observation of the quantum paradox of separation of a single photon from one of its properties James M. Ashby, Peter D. Schwarz, and Maximilian Schlosshauer Phys. Rev. A 94, 012102 – Published 1 July 2016**

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| (2906) Authors report an experimental realization of the quantum paradox of (e) the separation of a single photon from (e) one of its properties (the so-called "quantum Cheshire cat"). |
| (2907) Authors use a modified Sagnac interferometer with (e&eb) displaced paths to produce (eb) appropriately pre- and postselected states of (e) heralded single photons. |
| (2908) Weak measurements of photon presence and (e&eb) circular polarization are performed in each arm of (e) the interferometer by introducing (eb) weak absorbers and small polarization rotations and analyzing changes in (eb) the postselected signal. |
| (2909) The absorber is found to have (e) an appreciable effect only in one arm of the interferometer, while (e) the polarization rotation significantly affects (e&eb) the signal only when (e) performed in the other arm. |
| (2910) Authors carry out both sequential and simultaneous weak measurements and find (eb) good agreement between measured and predicted weak values. |
| (2911) In the language of Aharonov et al. and in the sense of the ensemble averages described by (e) weak values, the experiment establishes (eb) the separation of a particle from one its properties during the passage through (e&eb) the interferometer. |

**Viewpoint: Relativity Gets Thorough Vetting from LIGO Frans Pretorius, Department of Physics, Princeton University, Princeton, NJ 08544, USA May 31, 2016• Physics 9, 52 The gravitational-wave signal observed by the LIGO detectors shows no deviation from what general relativity predicts. Figure caption expand figure Frans Pretorius, APS/Carin Cain**

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| (2912) Figure 1: The signal from one of the LIGO detectors in Hanford, Washington, is shown with two representations of (e) the best-fit numerical relativity (NR) waveform. |
| (2913) The filtered NR waveform illustrates how the raw waveform is perceived by (e) the detector, |
| (2914) On September 14, 2015, the LIGO interferometers made the first direct measurement of (e) gravitational waves [1]. |
| (2915) The historic discovery came nearly 100 years after Einstein showed that his general relativity theory predicts (eb) the existence of gravitational waves and over 50 years since (e) the first quest to observe the Universe in gravitational waves [2]. |
| (2916) The observed event, labeled GW150914, was (=) the collision of two black holes, one weighing 29 solar masses, the other 36 solar masses. |

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| (2917) This binary coalesced into (e&eb) a remnant black hole weighing 62 solar masses and rotating (e&eb) at 67% the maximum rate allowed (eb) for a black hole. |
| (2918) The most statistically significant portion of (e) the signal came from (e) the last several gravitational-wave cycles of the collision, where (e) the most nonlinear, strong-field dynamics of the theory manifest. |
| (2919) One couldn't have asked for a better first event to learn about (e&eb) this regime of general relativity. |
| (2920) The first thing that GW150914 can teach us is how closely nature follows (eb) Einstein's prescriptions for (e&eb) colliding black holes. |
| (2921) Researchers from the LIGO team and the Virgo team have performed a series of tests to see if (e) any violations of (e) general relativity might be lurking in the data of GW150914 [3]. |
| (2922) They first determined (eb) the theoretical waveform, or "template," from (e) general relativity simulations that best fit the measured signal (see Fig. 1). |
| (2923) Their main finding is that the residual signal—obtained by (e) subtracting this template from (e) the data—is consistent with (e&eb) noise. |
| (2924) This is, in a sense, the "everything" test, as (e) the template folds in (e&eb) all that general relativity predicts about the merger: the decay of the last few orbits of (e) the two black holes, their violent collision (e&eb) and astonishingly rapid relaxation to (e) a quiescent rotating black hole, the propagation of (e) gravitational waves through (e&eb) the cosmos, and finally the transformation of (e) these waves into(e&eb) a chirp in the LIGO detectors. |
| (2925) Given the large signal-to-noise ratio of 24, the close agreement between the template and (=) the data puts limits on (e) alternative models. |
| (2926) Specifically, the predictions of these models can only disagree with (e&eb) the general relativity prediction for GW150914 by, at most, 4%. |
| (2927) In addition to this comprehensive test, the LIGO/Virgo teams also performed other checks that pick apart the signal, in order to address particular aspects of general relativity and black hole physics. In one such test, they compared the early and (e&eb) late parts of the waveform. |
| (2928) From the waves generated during (e&eb) the initial inspiral, one can infer (eb) the masses and spins of the binary. |
| (2929) Given these numbers, general relativity allows (eb) the mass and spin of (e) the remnant black hole to be calculated. |
| (2930) The researchers compared this result to (e&eb) an independent measurement of (e) the mass and spin obtained from (e) the late-time ringdown, and to within (eb) uncertainties, the two estimates were (=) consistent. |
| (2931) Another test isolated coefficients in the so-called Post Newtonian (PN) series describing (eb) the inspiral. |
| (2932) The PN series is (=) an expansion of the gravitational-wave signal in powers of (e) the velocity vv of the sources relative to (e&eb) the speed of light cc. |
| (2933) Observations of binary pulsars, which show (eb) a decrease in their orbital period due to (e) gravitational-wave emission, have constrained (e) PN terms, especially at low order. |
| (2934) For higher-order terms associated with (e&eb) strong-field emission, the LIGO/Virgo teams were able to improve (eb) the limits by many orders of magnitude compared to (e&eb) the binary pulsar constraints. |

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| (2935) At first glance it might seem these new results simply add to (e) the litany of previous tests that general relativity passed with flying colors. |
| (2936) However, GW150914 is vastly more significant than that. Other tests have shown that general relativity is valid in (eb)the weak-field regime, but GW150914 gives (eb) the first direct observational evidence that general relativity is the correct theory of nature governing (e&eb) the extreme spacetime dynamics around the event horizon of (e) a black hole. |
| (2937) One could argue that binary pulsar observations already told us that gravitational waves exist [4], but GW150914 is more luminous in gravitational waves than pulsar binaries by (e) at least 24 orders of magnitude. |
| (2938) It's true that the evolution of the Universe as a whole is driven by (e) strong gravity, but it would be difficult to argue that precision cosmological measurements can constitute (eb) a test of general relativity when (e) one of the dominant inputs, dark energy, is still a mystery. |
| (2939) As far as black holes are concerned, astronomers have (e) observed ultra compact dark objects that are consistent with (e&eb) black holes. |
| (2940) The most striking data comes from (e) stellar orbits about the 4 million solar mass objects in (eb) the center of our Galaxy [5], but present observations cannot constrain it to be (=) a black hole as described by general relativity. |
| (2941) GW150914 has given us the first morsels of quantifiable evidence that the rich physics of black holes is as predicted by general relativity. For example, the LIGO observations are consistent with (e&eb) so-called cosmic censorship [6], which says that gravitational singularities must remain hidden behind an (eb) event horizon. |
| (2942) If a "naked" singularity were produced in (eb) the energetic collision of GW150914, the data would likely have shown (eb) some sort of anomaly, but in exactly what form, we don't know. |
| (2943) Another example involves the "no-hair" theorem [7], which says that the only properties of a black hole are (=) its mass, spin, and charge (though astrophysical black holes are expected to have negligible charge). |
| (2944) By consequence, these properties should determine (eb) the frequency and decay constants of (e) all the infinitely many ringdown modes of (e) a black hole. |
| (2945) With GW150914, only the least-damped ringdown mode is (=) visible above the noise, but when combined with (e&eb) the remnant mass and spin inferred from (e) the inspiral, this is enough to count as (=) the first piece of evidence for a black hole's baldness. |
| (2946) GW150914 should also be a test for (e) alternative theories of gravity or exotic alternatives to (e) black holes within general relativity, such as boson stars [8]. |
| (2947) Some of these alternative models are concerned with (e&eb) explaining dark energy or dark matter and do not (e) therefore, predict any deviation from relativity at (eb) the scale of GW150914. |
| (2948) However, other models do in fact deviate from (e) conventional relativity in their predictions of (e) the gravitational-wave signal from (e) a binary merger. |
| (2949) The largest deviations should presumably occur during (e&eb) the last several gravitational-wave cycles, where (e) the signal is strong, and deviations should be (=) easy to detect or constrain. |
| (2950) Embarrassingly though, the majority of alternatives that claim deviations can't (e) really say much about this event, as the two-body problem hasn't been solved in these theories. |
| (2951) These models can only reproduce (eb) a fraction of the full gravitational-wave signal, (eb) so the strong constraints one might have hoped for (e) cannot be leveled. |

(2952) Exotic alternatives within (eb) general relativity are in a similar boat, or worse for (e) objects like gravastars, which have (e) no well-defined theoretical framework to say what happens when (e) two such objects collide.

(2953) Regardless, what is clear from the data is that any non-black-hole solution must imagine a form of (e) exotic matter with a viscosity larger than (e) any known material to explain (eb) the rapid damping of oscillations observed during the ringdown.

(2954) The inferred event rate for heavy black hole mergers like GW150194 is high, and LIGO has another factor of 3 improvement in sensitivity slated before it reaches (eb) design sensitivity.

(2955) Within 5 years then, we can expect dozens of detected events, with the loudest having a signal-to-noise ratio upwards of (e) 100.

(2956) This will allow (eb) for exquisite insight into the nature of extreme gravity beyond (e) what GW150194 has already gifted us with. **This research is published in Physical Review Letters**.

**Tests of General Relativity with GW150914 B. P. Abbott et al. (LIGO Scientific and Virgo Collaborations) Phys. Rev. Lett 116, 221101 – Published 31 May 2016**

(2957) The LIGO detection of GW150914 provides (eb) an unprecedented opportunity to study the two-body motion of (e) a compact-object binary in (eb) the large-velocity, highly nonlinear regime, and to witness the final merger of (e) the binary and the excitation of uniquely relativistic modes of (e) the gravitational field.

(2958) Authors carry out several investigations to determine whether GW150914 is consistent with (e&eb) a binary black-hole merger in general relativity.

(2959) They find that the final remnant's mass and spin, as determined from (e) the low-frequency (inspiral) and (e&eb) high-frequency (postinspiral) phases of the signal, are mutually consistent with (e&eb) the binary black-hole solution in general relativity.

(2960) Furthermore, the data following the peak of GW150914 are consistent with (e&eb) the least-damped quasinormal mode inferred from (e) the mass and spin of the remnant black hole.

(2961) By using waveform models that allow (eb) for parametrized general-relativity violations during (e&eb) the inspiral and merger phases, authors perform quantitative tests on (e&eb) the gravitational-wave phase in the dynamical regime and they determine the first empirical bounds on (e&eb) several high-order post-Newtonian coefficients.

(2962) Authors constrain (e) the graviton Compton wavelength, assuming (eb) that gravitons are dispersed in (eb) vacuum in the same way as particles with mass, obtaining a 90%-confidence lower bound of 1013 km.

(2963) In conclusion, within our statistical uncertainties, we find no evidence for violations of (e) general relativity in the genuinely strong-field regime of gravity. **Received 26 March 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.116.221101 © 2016 American Physical Society**

**Focus: Background Noise of Gravitational Waves March 31, 2016• Physics 9, 33**

(2964) Assuming LIGO's recently detected black hole merger is not unusual; researchers revised upward their estimate of (e) the strength of the background noise coming from (e) distant mergers across the Universe. Figure caption expand figure Caltech/MIT/LIGO Lab (deleted due to spatial constraints)

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| (2965) Catching some waves The LIGO-Livingston interferometer consists of (e) two 4-km-long arms where (e) laser light bounces back and forth. Interference between the lasers from the two arms can signal tiny changes in (e&eb) the lengths of the arms resulting from gravi... |
| (2966) The recent detection of gravitational waves emitted from (e) a pair of merging black holes implies (eb) the existence of a large population of similar pairs, which should be generating (eb) a background of waves too faint to be detected with (e&eb) present technology. |
| (2967) Using this information, researchers now estimate that this background could be (=) roughly 10 times stronger than previously expected and might be found by (e) sensitive instruments in only a few years. |
| (2968) The background signal could help astrophysicists better understand how black hole binaries form and also pave the way to detecting the waves emitted by (e) the expanding Universe shortly after the big bang. |
| (2969) Gravitational waves were detected for the first time on 14 September 2015 by the Laser Interferometer Gravitational-Wave Observatory (LIGO). |
| (2970) By reflecting laser light between mirrors set 4 km apart, the team was able to detect (e&eb) tiny fluctuations in the distance between the mirrors as (=) the waves passed. |
| (2971) The waves were triggered by (e) an event, called GW150914; involving (e&eb) the merger of a pair of black holes having (e) masses of about 29 and 36 times that of the Sun. |
| (2972) Researchers expect that similar pairs of black holes, as well as neutron stars, should be merging constantly, creating (eb) gravitational waves that should bombard (e&eb) the Earth from all directions and make the equivalent of (=) a faint popping noise in a detector. |
| (2973) Previous estimates of the likely strength of this background had to make assumptions about (e&eb) how many pairs there might be and about the distribution of their masses. |
| (2974) These estimates remained speculative, as no actual binaries were known. |
| (2975) "What the recent observation has changed," says Nelson Christensen of Carleton College in Minnesota and the LIGO Scientific Collaboration (LSC), "is that we now know there really is (=) a population of black hole binaries. We have a small bit of information about it." |
| (2976) The new analysis—carried out by LSC researchers working with (e&eb) the European Virgo collaboration—estimates (eb) the energy density in the gravitational wave background in (eb) the frequency range below 100 Hz, where (e) the LIGO detector is most sensitive. |
| (2977) To make this estimate, they assumed (eb) a population of black hole binaries of various sizes, having (e) a range of distances from Earth, and created through (e&eb) a variety of scenarios. |
| (2978) They then calculated the amount of gravitational radiation that would result from (e) the merging of these objects. |
| (2979) The team then adjusted this estimate by using (e) details of the GW150914 event—especially the masses of the black holes involved—to make (eb) a most likely guess at the actual distribution of black hole binaries. |
| (2980) The idea is that the observed merger is (=) most likely representative of (e) the general population, rather than being extremely unusual. |
| (2981) If the black hole binary population is represented as (=) a bell curve, or Gaussian, the event should lie within (eb) the main bulge, not in the tails. |
| (2982) The observed mass of GW150914 suggests (eb) that the main population of black hole binaries could be (=) roughly 20 times more massive than expected, which would generate (eb) a background roughly 10 times "louder" than (e) previous estimates. |

(2983) However, Christensen emphasizes that the result has (e) lots of uncertainty, as it's based on (e) only one event.

(2984) "This estimate seems very plausible to me," says astrophysicist Maura McLaughlin of the University of West Virginia in Morgantown. She is a member of the NANOGRAV collaboration, which tries to detect gravitational waves at much lower frequencies than LIGO can observe. "I can't see any holes in their arguments. Unless LIGO got really lucky, then there must be many more of these systems than previously thought." This level of gravitational background implies (eb) that the LIGO and Virgo detectors might be sensitive enough to see it in 2–5 years, sooner than physicists previously expected, says Christensen, which has led to much anticipation. "Everyone is super excited," he says.

(2985) Detecting the background would provide (eb) a confirmation of the presence of gravitational waves carrying (e&eb) energy throughout the Universe. It would also be useful, the researchers suggest, in gaining (e) further understanding of precisely how black holes form.

(2986) And subtracting this background might someday allow (eb) detection of the fainter waves coming from (e) the expansion of the Universe moments after the big bang. This research is published in Physical Review Letters. –Mark Buchanan Mark Buchanan is a freelance science writer in Normandy, France.

**GW150914: Implications for the Stochastic Gravitational-Wave Background from Binary Black Holes B. P. Abbott et al. (LIGO Scientific Collaboration and Virgo Collaboration) Phys. Rev. Lett. 116, 131102 – Published 31 March 2016**

(2987) The LIGO detection of the gravitational wave transient GW150914, from (e) the inspiral and merger of (e) two black holes with masses $\gtrsim 30 M_{\odot}$, suggests (eb) a population of binary black holes with (e&eb) relatively high mass.

(2988) This observation implies (eb) that the stochastic gravitational-wave background from binary black holes, created from (e) the incoherent superposition of (e) all the merging binaries in the Universe, could be (=) higher than previously expected.

(2989) Using the properties of GW150914, authors estimate the energy density of (e) such a background from (e) binary black holes.

(2990) In the most sensitive part of the Advanced LIGO and Advanced Virgo band for (e) stochastic backgrounds (near 25 Hz), authors predict $\Omega_{GW}(f=25\ \mathrm{Hz}) = 1.1^{+2.7}_{-0.9} \times 10^{-9}$ with 90% confidence.

(2991) This prediction is robustly demonstrated for (e) a variety of formation scenarios with (e&eb) different parameters.

(2992) The differences between models are (=) small compared to the statistical uncertainty arising from (e) the currently poorly constrained (e) local coalescence rate.

(2993) Authors conclude that this background is potentially measurable by (e) the Advanced LIGO and Advanced Virgo detectors operating at their projected final sensitivity. DOI:http://dx.doi.org/10.1103/PhysRevLett.116.131102 © 2016 American Physical Society

**Study opens new realms of light-matter interaction July 15, 2016 by David L. Chandler. Some sentences might have been omitted due to spatial constraints. Please bear with me.**

(2994) Emission spectra are a widely used method for (e) identifying chemical compounds; the bright lines reveal (eb) the different frequencies of light that can be emitted by (e) an atom.

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| (2995) Here, a normal emission spectrum for (e) an atom in a high-energy state (top) is compared to (e&eb) the emission from the same atom placed just a few nanometers (billionths of a meter) away from (e) graphene that has been doped with (e&eb) charge carriers (bottom). |
| (2996) For each energy-level transition, an orange line (or purple cloud) appears if (e) that transition is estimated to be faster than (e) one per microsecond — making it frequent enough to be (=) observed. **Credit: Massachusetts Institute of Technology** |
| (2997) A new MIT study could open up new areas of technology based on (e) types of light emission that had been thought to be (=) "forbidden," or at least so unlikely as to be (=) practically unattainable. |
| (2998) The new approach, the researchers say, could cause (eb) certain kinds of interactions between light and (e&eb) matter, which would normally take billions of years to happen, to take place instead (e&eb) within billionths of a second, under (e&eb) certain special conditions. The findings, based on a theoretical analysis, are reported today in the journal Science in a paper by MIT doctoral student Nicholas Rivera, Department of Physics Professor Marin Soljačić, Francis Wright Davis Professor of Physics John Joannopoulos, and postdocs Ido Kaminer and Bo Zhen. |
| (2999) Interactions between light and matter, described by (e) the laws of quantum electrodynamics, are the basis of a (e) wide range of technologies, including (e) lasers, LEDs, and atomic clocks. |
| (3000) But from a theoretical standpoint, "Most light-matter interaction processes are 'forbidden' by (e) electronic selection rules, which limits (e) the number of transitions between energy levels we have access to," Soljačić explains. |
| (3001) For example, spectrograms, which are used to (e) analyze the elemental composition of materials, show (eb) a few bright lines against (e&eb) a mostly dark background. |
| (3002) The bright lines represent (eb) the specific "allowed" energy level transitions in (eb) the atoms of that element that can be accompanied by (e) the release of a photon (a particle of light). |
| (3003) In the dark regions, which make up most of the spectrum, emission at those energy levels is (=) "forbidden." |
| (3004) With this new study, Kaminer says, "we demonstrate theoretically that these constraints can be lifted" using (e) confined waves within (e&eb) atomically thin, 2-D materials. |
| (3005) "We show that some of the transitions which normally take the age of the universe to happen could be made to happen within (eb) nanoseconds. |
| (3006) Because of this, many of the dark regions of a spectrogram become (=) bright once an atom is (=) placed near a 2-D material." |
| (3007) Electrons in an atom have (e) discrete energy levels, and when they hop from one level to another they give off (eb) a photon of light, a process called spontaneous emission. |
| (3008) But the atom itself is much smaller than the wavelength of the light that gets emitted—about 1/1,000 to 1/10,000 as big—substantially impairing (e&eb) the interactions between the two. |
| (3009) The trick is, in effect, to "shrink" (e) the light so it better matches the scale of the atom, as (=) the researchers show in their study. |
| (3010) The key to enabling a whole range of interactions, specifically transitions in (eb) atomic states that relate to absorbing or emitting light, is (=) the use of a two-dimensional material called graphene, in which light can interact with (e&eb) matter in the form of plasmons, a type of electromagnetic oscillation in the material. |
| (3011) These plasmons, which resemble photons but have wavelengths hundreds of times shorter, are (=) very narrowly confined in the graphene, in a way that makes some kinds of interactions with (e&eb) that matter many orders of magnitude more likely than (e) they would be in ordinary materials. |

(3012) This enables a variety of phenomena normally considered unattainable, such as (e) the simultaneous emission of multiple plasmons, or two-step light-emitting transitions between (e&eb) energy levels, the team says.

(3013) This method can enable the simultaneous emission of (e) two photons that are "entangled," meaning (eb) they share the same quantum state even when (e) separated.

(3014) Such generation of entangled photons is (=) an important element in quantum devices, such as those that might be used for cryptography.

(3015) Making use of these forbidden transitions could open up (eb) the ability to tailor the optical properties of materials in ways that had not been thought possible, Rivera says.

(3016) "By altering these rules" about the relationship between light and (e&eb) matter, "it can open (eb) new doors to reshaping (e&eb) the optical properties of materials."

(3017) Kaminer predicts that this work "will serve as (=) a founding piece for (e) the next generation of studies on light-matter interactions" and could lead to (eb) "further theoretical and experimental advances in many fields which rely on (e) light-matter interactions, including (e) atomic, molecular and optical physics, photonics, chemistry, optoelectronics, and many others."

(3018) Beyond its scientific implications, he says, "this study has possible applications across multiple disciplines, since (e) in principle it has potential to enable the full use of (e) the periodic table for optical applications."

(3019) This could potentially lead to (eb) applications in spectroscopy and sensing devices, ultrathin solar cells, and new kinds of materials to absorb solar energy, organic LEDs with higher efficiencies, and photon sources for (e) possible quantum computing devices."From the standpoint of fundamental science, this work lays the groundwork for a subfield that just a few years ago was difficult to imagine and until now was largely unexplored," Soljačić says.

(3020) "Two-dimensional materials confine (e) fields to a surface and motion to a plane, making possible (eb) many effects that are orders-of-magnitude too weak to (e) appear in a bulk volume," says Jason Fleischer, an associate professor of electrical engineering at Princeton University, who was not involved in this research.

(3021) This work, he says, "systematically explores how 2-D materials improve light-matter interactions, laying a theoretical foundation for faster electronic transitions, enhanced sensing, and better emission, including the compact generation of broadband and quantum light." More information: N. Rivera et al: Shrinking light to allow (eb) forbidden transitions on the atomic scale, Science (2016). DOI: 10.1126/science.aaf6308 Journal reference: Science search and more info website: Provided by: Massachusetts Institute of Technology search and more info website. This story is republished courtesy of MIT News (web.mit.edu/news office/), a popular site that covers news about MIT research, innovation and teaching.

**Viewpoint: Peering Back to the Beginning of Time Lawrence M. Krauss, Department of Physics and School of Earth and Space Exploration, Arizona State University, Tempe, AZ 85287, USA and Mount Stromlo Observatory, Research School of Astronomy and Astrophysics, Australian National University, Weston, ACT 2611, Australia June 19, 2014• Physics 7, 64**

(3022) The BICEP2 collaboration reports the detection of BB-mode polarization of (e) the cosmic microwave background—a signal that might originate from (e) gravitational waves created by (e) inflation during the very earliest moments in the evolution of our Universe. Figure caption expand figure P. A. R. Ade et al. [1]; Image on homepage: BICEP2 Collaboration (deleted due to restrictions on space)

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| (3023) Figure 1: (Top) BB-mode polarization pattern observed by (e) the BICEP2 telescope [1]. |
| (3024) The twisting pattern in the polarization of (e) the cosmic microwave background (CMB) might originate from (e) gravitational waves created by inflation in the very first moments of the universe. |
| (3025) Inflation—the hypothesis that the Universe underwent (e&eb) a phase of superluminal expansion in a brief period following the big bang—has (e) the potential of explaining, from first principles, why the Universe has (e) the structure we see today. |
| (3026) It could also solve outstanding puzzles of standard big-bang cosmology, such as why the Universe is, to a very good approximation, (=) flat and isotropic (i.e., it looks the same in all directions). |
| (3027) Yet we do not yet have a compelling model, based on (e) fundamental particle physics principles, that explains inflation. |
| (3028) And despite its explanatory power and a great deal of suggestive evidence, we still lack an unambiguous and direct probe of (e) inflation. |
| (3029) Theorists have developed different models for inflation, which all share (e&eb) a common, robust prediction: Inflation would have created (eb) a background of gravitational waves that could have (eb) an observable effect. |
| (3030) These waves would cause (eb) subtle, characteristic distortions of (e) the cosmic microwave background (CMB)—the oldest light in the Universe, released when (e) photons decoupled from matter and the Universe became (=) transparent to radiation |
| (3031) Until now, such a background of gravitational waves has (e) escaped detection. It is therefore very exciting that the BICEP2 (Background Imaging of Cosmic Extragalactic Polarization) collaboration has now reported in Physical Review Letters the observation of polarization modes of CMB radiation [1] that are consistent with (e&eb) such a background. |
| (3032) If confirmed, BICEP2's finding would provide (eb) the first direct probe of physical processes that occurred a mere 10−3610-36 seconds after (e) the big bang. |
| (3033) The result also offers a rare opportunity to directly test theoretical ideas, including inflation and grand unification (the theory that merges the electromagnetic, weak, and strong forces into a single force), at energies 13 orders of magnitude greater than (e) we can currently access at the Large Hadron Collider. |
| (3034) Inflation is based on (e) the possibility that a scalar field in the early Universe got momentarily stuck in (eb) a "false vacuum state"—a metastable condition in which the Universe wasn't at (eb) its global potential minimum. |
| (3035) The energy stored in this vacuum state could drive (eb) an exponential expansion of spacetime (similar to, but much faster than, the expansion due to dark energy seen today) before (e) a phase transition to the true vacuum state occurs, releasing (eb) this energy into (e&eb) particles and radiation. |
| (3036) This creates (eb) an initially hot state for (e) a subsequent big-bang-like expansion from (e) which our Universe subsequently evolved [2]. |
| (3037) If the Universe grew by at least a factor of 10261026, (eb) inflation can explain (eb) both the apparent flatness and isotropy of the Universe [2]: |
| (3038) The expansion would stretch (eb+ eb-dot) any initial curvature of (e) the Universe to near flatness, and smooth out (e&eb) any initial inhomogeneities, just as stretching (eb+ ebdot) a piece of fabric would do. |
| (3039) Isotropy is justified by the fact that the Universe, before inflation, was once much smaller and could have been contained in (eb) a causally connected region that was in thermal equilibrium, resulting in (eb) the approximately uniform, isotropic temperature distribution observed on (eb) large scales today. |

(3040) Importantly, the theory can explain the (eb) origin of structure in the Universe: Microscopic quantum fluctuations associated with (e&eb) the scalar field driving inflation could have been "frozen in" during the exponential expansion, resulting, (eb) after inflation, in density fluctuations that might have seeded all observed large-scale structures in the Universe, including (e) all visible galaxies and clusters of galaxies.

(3041) The signs of such primordial fluctuations have indeed been detected through (e&eb) their effect on (e&eb) the CMB, first by the Cosmic Background Explorer (COBE) [3] and later by other missions.

(3042) These observational studies detected small temperature variations (anisotropies) in (eb) the CMB and found a spectrum compatible with (e&eb) that predicted by inflation.

(3043) The magnitude of these fluctuations suggests (eb) an energy scale of inflation (loosely speaking, the energy stored throughout space in (eb) the scalar field potential that drove inflation) close to (e) 10161016 giga-electron-volts (GeV).

(3044) But while consistent with inflation, such fluctuations are not (e) conclusive evidence: Their predicted shape and magnitude is known to vary with (e&eb) the details of the inflationary potential (whose form is not yet experimentally constrained).

(3045) Hence many different observations could be accommodated, and any agreement between theory and observation is at best suggestive. There is also a possibility that other, noninflationary sources might contribute. However, gravitational waves—the propagating ripples in spacetime predicted by Einstein in 1916—might provide (eb) direct evidence of inflation.

(3046) During inflation, gravitational fields would also fluctuate because of (e) quantum effects.

(3047) These fluctuations would similarly get "frozen in" and re-emerge after inflation as (=) classical gravitational waves, whose amplitude depends only on (=) the energy stored in the field driving (eb) inflation, and not on other model-dependent details [4].

(3048) If the scale of inflation were indeed as high as a few times 10161016 GeV, this background could produce (eb) specific temperature anisotropies in the CMB [5].

(3049) These anisotropies have thus far not been detected by (e) the most recent searches by the WMAP and PLANCK satellites, and their absence was used to (e) put upper limits on the scale of inflation somewhat below 10161016 GeV.

(3050) But a more sensitive CMB probe of (e) gravitational waves exists.

(3051) Large-wavelength gravitational waves produced by (e) inflation would have alternatively compressed and expanded (e&eb) space, producing (eb) spatial variations in (eb) the intensity of radiation bombarding free electrons in (eb)the early Universe (see, e.g., Ref: [6]).

(3052) As photons decoupled from (e) matter and CMB radiation was released, scattering from (e) the undulating sea of electrons would have created (eb) patterns in the polarization of CMB.

(3053) Such polarization anisotropies are predicted in (eb) inflationary models to be (=) most visible when (e) looking at the sky on spatial scales comparable to (e&eb) or somewhat larger than the size of (e) the Universe's horizon at (eb) the time when the CMB was created.

(3054) This corresponds to (e&eb) about 1 degree of the field of view of the sky.

(3055) Over the past decade, various CMB polarization experiments have thus chased the observation of polarization anisotropies on (e&eb) these scales but, so far, also yielded (eb) merely upper limits on a possible gravitational background.

(3056) All of this has now changed with (e&eb) the BICEP2 collaboration's [1] report of (e) a signal that has the predicted (eb) characteristics of such a background at precisely the angular scales where (e) the inflationary signal is expected (see Fig. 1, top deleted due to spatial restraints).

(3057) The BICEP2 Collaboration uses (e) a microwave telescope located at the South Pole, with (e&eb) an array of polarization-sensitive detectors pointing at (eb) a small field (a solid angle of about 11 percent of the full sky) with a resolution of approximately 0.50.5 degrees, thus sufficient to probe the relevant angular scales for the detection of (e) polarization anisotropies originating from (e) **inflation**.

(3058) BICEP2's search is based on (e) two expected features of primordial gravitational waves.

(3059) First, gravitational waves would produce (eb) a characteristic spatial polarization pattern on CMB—a consequence of (e) their quadrupole nature.

(3060) One can decompose (e&eb) any spatial polarization pattern into (e&eb) two orthogonal components, a "curl-free" component and (e&eb) a component involving a twisting pattern.

(3061) In analogy to electromagnetic modes, these polarization patterns are called "EE modes" and (e&eb) "BB modes," respectively (see Fig. 1).

(3062) Scalar fluctuations in the density of (e) matter would produce (eb) EE modes, while (e) gravitational waves would instead generate (eb) a BB-mode pattern.

(3063) The second feature is the aforementioned characteristic spatial scale of (e) expected CMB polarization signatures, including (e) wavelength scales comparable to (e&eb) or larger than the horizon scale at the time the CMB was created.

(3064) Since inflation would produce (eb) both gravitational waves and (e&eb) scalar density perturbations, BICEP2, like all CMB-polarization experiments, quote their results in terms of (e&eb) the ratio rr of the amplitude of (e) tensor perturbations (associated with (e&eb) gravitational waves) to (e) the amplitude associated with (e&eb) scalar density fluctuations (reflected in (eb) the magnitude of CMB temperature fluctuations measured by (e) experiments like COBE).

(3065) A nonzero rr implies (eb) the detection of a BB-mode pattern and primarily measures (eb) the strength of any possible gravitational wave signal which, if (e) it comes from inflation, is then determined by (e) its energy scale (e.g., see Ref. [7] for a review).

(3066) BICEP2 reported a nonzero value of r=0.2(+0.07−0.05)r=0.2(+0.07-0.05), on scales ranging from (e) the horizon scale up to 4–5 times larger, with (e&eb) a claimed significance of better than 5σ5σ.

(3067) If confirmed, a value of 0.2 corresponds to (e&eb) an energy density during (e&eb) inflation of ∼1.2×1016∼1.2×1016 GeV.

(3068) The claimed signal is far larger than expected, partly because (e) the PLANCK satellite, despite not having yet reported its polarization measurements, has (e) previously put an upper limit on (eb) rr (<0.11<0.11). PLANCK's estimate is based on (e) the absence of a gravitational wave signature on (eb) the CMB temperature anisotropy on (eb) its largest scales (PLANCK observes the whole sky and therefore angular scales up to 360360 degrees).

(3069) The apparent discrepancy between the PLANCK upper bound and (e&eb) the BICEP2 detection might be a cause for concern.

(3070) However, the two bounds probe different spatial scales, and the BICEP collaboration has noted that, while tensor perturbations are expected to be (=) roughly scale independent, some scale dependence of (e) scalar density perturbations could relax (e&eb) the PLANCK constraint and bring (eb) the two results into agreement.

(3071) Of greater concern, of course, is the possibility that the signal could be determined by (e) foreground contamination—any source along (e&eb) the line of sight of the detector that could distort or contaminate (e&eb) the primordial CMB signal.

(3072) BICEP2 tested for (e) a wide variety of possible foregrounds, claiming (eb) they cannot be responsible for (e) the observed signal.

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| (3073) One obvious candidate arises from (e) gravitational lensing: The intervening mass of (e) astrophysical bodies could distort (e&eb) the spatial pattern of polarization in a way that that could transform (e&eb) an EE-mode signal into (e&eb) a BB-mode signal. |
| (3074) The BICEP collaboration has analyzed this possibility extensively. Based on (e) measurements of (e) CMB temperature (and hence density perturbations) and on (e&eb) models of the mass distribution, they argue that such effects would only be (=) significant on smaller angular scales than those contributing to their measurements of rr. |
| (3075) But since BICEP2's first announcement on March 17, 2014, a number of scientists have raised another foreground concern due to (e) the possible impact of (e&eb) our Galaxy's dust. |
| (3076) BICEP2 modeled dust contamination scenarios based on (e) a variety of theoretical models, using data from (e) its predecessor BICEP1 and preliminary dust estimates from (e) the PLANCK satellite. |
| (3077) They concluded that dust could not (e) reproduce the magnitude of the observed signal. |
| (3078) However, PLANCK has in the meanwhile released (eb) more complete measurements of variation in polarized dust foregrounds in (eb) many regions of the Galaxy. |
| (3079) Their analysis increases the uncertainty associated with (e&eb) the dust contribution to (e) the BICEP2 result and also accommodates (e) a potentially larger polarized dust contribution. |
| (3080) Some researchers argued that dust polarization at a level 2–3 times higher than that assumed by BICEP2 is possible and could therefore plausibly be responsible for (e) the entire signal (e.g., see Ref: [8]). |
| (3081) While the BICEP2 collaboration continues to maintain that this is extremely unlikely based on (e) their estimates, many other researchers are more skeptical. |
| (3082) In any case, a variety of upcoming measurements, from PLANCK and other detectors are (=) expected to resolve this important ambiguity within the coming year. |
| (3083) Even if the BICEP2 signal does arise from (e) primordial gravitational waves, various consistency checks will be required (e) before it can be definitively argued that the signal arises from (e) inflation. |
| (3084) Inflation is certainly the best-motivated possibility, and the spectral shape of the signal, at least for (e) the five largest-scale modes observed by BICEP2, agrees extremely well with (=) inflationary predictions (see Fig. 1, bottom). |
| (3085) Several results have already appeared that strongly constrain (e) other possible sources, ruling out (e) for example, the possibility of gravitational waves generated by (e) other possible phase transitions after inflation [9] |
| (3086) . This conclusion is reinforced by more careful analyses of the scale dependence of (e) the predicted signal (e.g., Ref. [10]). |
| (3087) But it will nevertheless be important to assure that all possible alternative sources of (e) gravitational waves are ruled out. |
| (3088) If it can be established that the BICEP2 represents (eb) a gravitational wave signal from (e) inflation, there is one additional exciting implication. |
| (3089) Standard calculations suggest (eb) that the gravitational wave signal from inflation is due to (e) quantum fluctuations in the gravitational field, and this appears to require (e) the existence of gravitons—the hypothetical elementary particles that mediate (e&eb) the force of gravity. |
| (3090) Research by Frank Wilczek and myself, based on (e) a dimensional analysis argument [11], suggests that,(eb) independently of the method used to calculate the gravitational wave spectrum, a quantum gravitational origin is (=) required. |
| (3091) If so, the BICEP2 result then implies (eb) that gravity is ultimately a quantum theory, a result of fundamental importance for physics. |

(3092) It is hard to imagine a more important or wider-ranging cosmological observation. If confirmed by upcoming experiments, the BICEP2 results will push the frontiers of empirical physics forward in an unprecedented way, offering (eb) the first direct observation of gravitational waves and of quantum gravity effects as well as a new probe of the earliest moments in the Universe's life and of physics at grand-unified energy scales. This will thus constrain both our understanding of fundamental interactions and the evolution of our Universe as a whole.

**Detection of B-Mode Polarization at Degree Angular Scales by BICEP2 P. A. R. Ade et al. (BICEP2 Collaboration) Phys. Rev. Lett 112, 241101 – Published 19 June 2014**

(3093) Authors report results from the BICEP2 experiment, a cosmic microwave background (CMB) polarimeter specifically designed to search for (e) the signal of inflationary gravitational waves in (eb) the B-mode power spectrum around $\ell \sim 80$.

(3094) The telescope comprised a 26 cm aperture all-cold refracting optical system equipped with (e&eb) a focal plane of 512 antenna coupled transition edge sensor 150 GHz bolometers each with temperature sensitivity of $\approx 300\ \mu K_{CMB}\sqrt{s}$. BICEP2 observed from (e) the South Pole for three seasons from 2010 to 2012.

(3095) A low-foreground region of sky with (e&eb) an effective area of 380 square deg was observed to (e) a depth of 87 nK deg in Stokes Q and U.

(3096) In this paper authors describe the observations, data reduction, maps, simulations, and results. They find an excess of B-mode power over the base lensed-ΛCDM expectation in the range $30<\ell<150$, inconsistent with (e&eb) the null hypothesis at (eb) a significance of $>5\sigma$.

(3097) Through jackknife tests and simulations based on (e) detailed calibration measurements authors show that systematic contamination is (=) much smaller than the observed excess.

(3098) Cross correlating against WMAP 23 GHz maps we find that Galactic synchrotron makes a negligible contribution to the observed signal. We also examine a number of available models of polarized dust emission and find that at their default parameter values they predict power ~(5–10)× smaller than the observed excess signal (with no significant cross-correlation with our maps). However, these models are not sufficiently constrained by external public data to exclude the possibility of dust emission bright enough to explain the entire excess signal. Cross correlating BICEP2 against 100 GHz maps from the BICEP1 experiment, the excess signal is confirmed with $3\sigma$ significance and its spectral index is found to be consistent with that of the CMB, disfavoring dust at $1.7\sigma$. The observed B-mode power spectrum is well fit by a lensed-ΛCDM+tensor theoretical model with tensor-to-scalar ratio $r=0.20+0.07-0.05$, with $r=0$ disfavored at $7.0\sigma$. Accounting for the contribution of foreground, dust will shift this value downward by an amount which will be better constrained with upcoming data sets. Received 4 April 2014 DOI:http://dx.doi.org/10.1103/PhysRevLett.112.241101 This article is available under the terms of the Creative Commons Attribution 3.0 License. Further distribution of this work must maintain attribution to the author(s) and the published article's title, journal citation, and DOI. Published by the American Physical Society

**(3099) Physicist claims to have observed quantum effects of Hawking radiation in the lab for the first time August 16, 2016 by Bob Yirka report**

(3100) This computer-simulated image shows a supermassive black hole at the core of a galaxy. The black region in the center represents the black hole's event horizon, where no light can escape the massive object's gravitational grip. The black …more

(3101) (Phys.org)—Jeff Steinhauer, a physicist at the Israel Institute of Technology, has published a paper in the journal Nature Physics describing experiments in which he attempted to create a virtual black hole in the lab in order to prove that Stephen Hawking's theory of radiation emanating from black holes is correct —though his experiments are based on sound, rather than light. In his paper, he claims to have observed the quantum effects of Hawking radiation in his lab as part of a virtual black hole—which, if proven to be true, will be the first time it has ever been achieved.

(3102) For many years, scientists believed that nothing could ever escape from a black hole. But in 1974, Stephen Hawking published a paper suggesting that something could—particles that are now called Hawking radiation. His idea was that if a particle (and its antimatter mate) appeared spontaneously at the edge of a black hole, one of the pair might be pulled into the black hole while the other escaped, taking some of the energy from the black hole with it—which would explain why black holes grow smaller and eventually disappear. Because such emissions are so feeble, no one has been able to measure Hawking radiation, so researchers have instead tried to build virtual black holes in labs to test the theory. One type of virtual black hole was proposed back in 1981 by Bill Unruh with the University of British Columbia—he suggested that an analogue might be created using water instead of light. He imagined a phonon existing at the edge of a waterfall—as the water speeds up, it begins to move faster than the speed of sound, causing it to be trapped. But if the phonon had an entangled mate that eluded the fall by moving away before getting caught up, it could escape. In this new effort, Steinhauer has built a device based on that idea and in so doing, claims he has observed an analogue of Hawking radiation.

(3103) The experiment consisted of creating an entangled pair of phonons sitting inside a bit of liquid that had been forced (via laser) to move very fast and then observing the action as one of the pair was pulled away as the liquid began to move faster than the speed of sound, while the other escaped—the fluid was a Bose-Einstein condensate of rubidium-87 atoms. After repeating the experiment 4,600 times Steinhauer became convinced that the particles were entangled, a necessity for a Hawking radiation analogue. His findings do not prove Hawking's theory to be true, of course, but they do appear to add a degree of credence that other researchers have thus far not been able to achieve. More information: Jeff Steinhauer. Observation of quantum Hawking radiation and its entanglement in an analogue black hole, Nature Physics (2016). DOI: 10.1038/nphys3863

**(3104) Abstract**

We observe spontaneous Hawking radiation, stimulated by quantum vacuum fluctuations, emanating from an analogue black hole in an atomic Bose–Einstein condensate. Correlations are observed between the Hawking particles outside the black hole and the partner particles inside. These correlations indicate an approximately thermal distribution of Hawking radiation. We find that the high-energy pairs are entangled, while the low-energy pairs are not, within the reasonable assumption that excitations with different frequencies are not correlated. The entanglement verifies the quantum nature of the Hawking radiation. The results are consistent with a driven oscillation experiment and a numerical simulation. Journal reference: Nature Physics search and more info website

Read more at: http://phys.org/news/2016-08-physicist-quantum-effects-hawking-lab.html#jCp

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1975); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as

the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 1985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_7}{dt} = [N(t)]_7[(F_7)(R(t)]_7 - [M]_7[R(t)]_7$$

$$\frac{d[R(t)]_7}{dt} = [B]_7\,[(R)_7(t - T_7)] - [D]_7[(R(t)]_7 - [N(t)]_7\left[G\left(R(t)\right)_7\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_7}{dt} = a_7[N(t)]_7 - a_7[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_7 = a_3[N(t)]_3 - a_4[N(t)]_7$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_7 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

These equations can be written for

**Viewpoint: Many-Body Localization Needs a Bath David A. Huse, Physics Department, Princeton University, Princeton, NJ 08544, USA and Institute for Advanced Study, Princeton, NJ 08540, USA July 5, 2016• Physics 9, 76 Theoretical work proves that interacting quantum systems can enter a many-body localized phase in which they cannot reach thermal equilibrium without an external bath. Figure caption expand figure Chris Laumann/University of Washington (deleted due to spatial restrictions). For references please see the original article**

(22) One striking feature of such MBL systems is that they have (e) strictly zero thermal and electrical conductivity: they are perfect insulators.

(23) However, even if "true to all orders" in perturbation theory, the conclusion by Basko, Aleiner, and Altshuler does not (e) rule out another possibility: many-body localization might only be (=) metastable, with (e&eb) the system appearing localized but actually thermalizing on (e&eb) long time scales because of (e) subtle effects that are not captured by (e) perturbation theory. One such scenario has been published earlier this year [8].

(24) Imbrie's new work [2] makes important progress towards a conclusive proof of a MBL phase by essentially ruling out (e) all such nonperturbative effects for certain systems.

Treating a particular model of interacting spins, he shows (eb) that all eigenstates of its Hamiltonian are (=) MBL.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.7b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 7a is at the

**TIME DELAY**

FIG. 7. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period,

which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.7$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types system takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_7}{dt} = N_7(t - T_7)F_7(R_7(t - T_7) - N_7(t)M_7((R_7(t) \qquad (5)$$

$$\frac{dR_7(t)}{dt} = B_7\big(R_7(t) - D_7(R_7(t)\big) - N_7(t)G_7(R_7(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_7}{dR_7}+\frac{N_7 dG_7}{dR_7}-\frac{dB_7}{dR_7}>\frac{N_7 G_7}{2M_7}\left[\frac{dF_7}{dR_7}+\frac{dM_7}{dR_7}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_7/dR_7$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_7$ This term will be defined by $P$.Further, we define two terms, also normalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_7$) and the other derived from the non-delayed or "immediate" processes($I_7$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_7=\frac{N_7 G_7}{2M_7}\frac{d(7nF_7)}{dR_7}(\geq 0) \quad (7a)$$

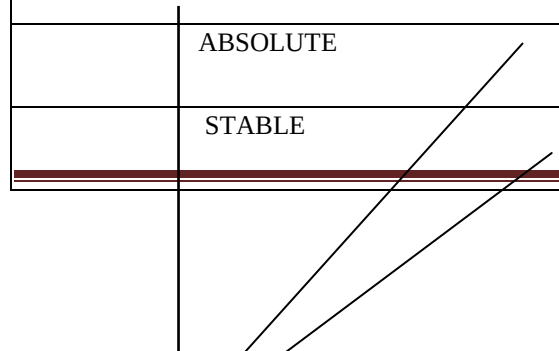
$$I_7=-\frac{N_7 G_7}{2M_7}\frac{d(7nM_7}{dR_7}(\geq 0) \quad (7b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_7>L_7-I_7$, must be modified ahe $dM_7/dR_7$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

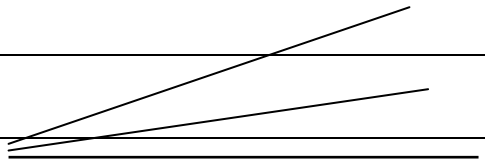
$$P_7>L_7-I_7, \quad \text{when} \quad P_7^2/4>I \quad (8a)$$

$$P_7>L_7-I_7+\frac{(I_7-P_7^2}{(L_7-I)} \quad \text{when} \quad P_7^2/4<I. \quad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,





**DELAYED PREDATOR CREDIT TYPES EFFECT WITH ITS WIDE RANGING RAMIFICATION**

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_7$; condition (8b) ensures that the total debit regulation, $P_7$, must always be positive even when $I_7$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_7$ and $L_7$. As $I_7$ increases relative to $L_7$ the zone of stability is enlarged; however,even when $I_7$is zero a significant zone of stability exists.

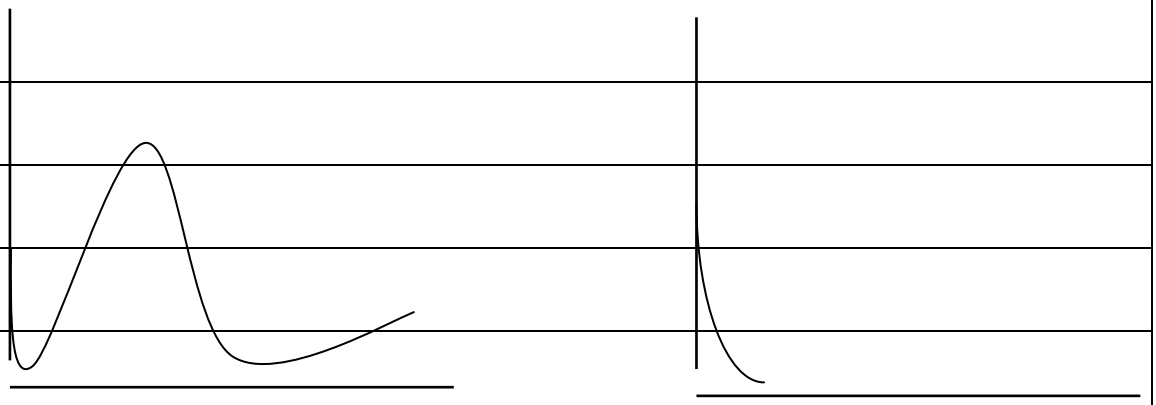
It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay

and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_7$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



**TIME DELAY**

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (7). Parameters in each case M=7; $dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7 = 0.05$ For (a) $NGdF/dR = 0.3, -N_7G_7dM_7/dR_7 = 0.6$; for (b)$NG_7dM_7/DR_7 = 0$

To scores of problems this relationship is important. Of them some of the examples are as follows:

**New look at the galaxy-shaping power of black holes**

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| (897) | Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of the plasma are moving. |
| (898) | "These giant black holes are (=) among the universe's most efficient energy generators, a hundred times more efficient than a nuclear reactor," said McNamara from Waterloo's Department of Physics and Astronomy in the Faculty of Science. |
| (899) | "Matter falling into the black hole is ripped apart, releasing (eb) vast amounts of energy in the form of high speed particles and thermal energy." |
| (900) | This heat is released from (e) just outside the black hole's event horizon (=) the boundary of no return. |
| (901) | The remaining matter gets absorbed into the black hole, adding to (e) its mass. |
| (902) | The released energy heats up (eb) the surrounding gas, creating (eb) bubbles of hot plasma that ripple through (e&eb) the cluster, just as bubbles of air rise up in (eb) a glass of champagne. |
| (903) | The research is shedding light on the crucial role that this hot plasma plays (e&eb) in galactic evolution. |
| (904) | Researchers are now tackling the foremost issue in the formation of (e) structure in the universe and asking: why don't most of the gas cool down, and form (eb) stars and galaxies? The answer seems to be that bubbles created by blasts of energy from (e) the black holes keep temperatures too high for (e) such structures to form. |
| (905) | "Any time a little bit of gas falls into the black hole, it releases (eb) an enormous amount of energy," said McNamara. |
| (906) | "It creates (eb) these bubbles, and the bubbles keep the plasma hot. |
| (907) | That's what prevents galaxies from (e) becoming even bigger than they are now." |
| (908) | Because plasma is invisible to the eye, and to optical telescopes, it wasn't until the advent of X-ray astronomy that the full picture began to emerge. In visible light, the Perseus cluster appears to contain (e) many individual galaxies, separated by (e) seemingly-empty space. |
| (909) | In an X-ray image, however, the individual galaxies are invisible, and the plasma atmosphere, centered on the cluster's largest galaxy, known as NGC 1275, dominates (e) the scene. |
| (910) | Although the black hole at the heart of NGC 1275 has (e) only one-thousandth of the mass of its host galaxy, and has (e) a much smaller volume, it seems to have (e) a huge influence on (e&eb) how the galaxy and how the surrounding hot plasma atmosphere evolve. |
| (911) | "It's as though the galaxy somehow knows (e) about this black hole sitting at the centre," said McNamara. |
| (912) | "It's like nature's thermostat, which keeps these galaxies from (e) growing. |
| (913) | If the galaxy tries to grow too fast, (eb) matter falls into the black hole, releasing (eb) an enormous amount of energy, which (eb) drives out the matter and prevents (e) it from forming new stars." |

**A new look at the galaxy-shaping power of black holes**

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| (914) | An X-ray image reveals (eb) the hot plasma that envelope (e) the Perseus cluster. Located some 240 million light years from earth, the |
| (915) | McNamara notes that the actual event horizon of the black hole is (=) about the same size as our solar system, making it as small compared to (e) its host galaxy as a grape is to the Earth. "What's going on in this tiny region is affecting a (e&eb) vast volume of space," he said. |

(916) Thanks to the black hole's regulatory effect, (eb) the gas that would have formed (eb) new stars instead remains hot plasma - whose properties Hitomi was designed to measure.

**Doomed satellite missions**

(917) Hitomi employed an X-ray spectrometer which measures (eb) the Doppler shifts in emissions from (e) the plasma; those shifts can then be used to (e) calculate the speed at which different parts of (e) the plasma are moving.

(918) At the heart of the spectrometer is a microcalorimeter; cooled to just one-twentieth of a degree above absolute zero, the device records (eb) the precise energy of each incoming X-ray photon.

(919) Getting an X-ray satellite equipped with a microcalorimeter into space has proved daunting: McNamara was deeply involved with NASA's Chandra X-ray Observatory, launched in 1999, that was initially set to include a microcalorimeter, but the project was scaled back due to budget constraints, and the calorimeter was dropped. Another mission with the Japanese space agency known as ASTRO-E was equipped with a microcalorimeter; it was set for launch in 2000, but the rocket exploded shortly after liftoff. A third effort, Japan's Suzaku satellite, launched in 2005, but a leak in the cooling system destroyed the calorimeter. Hitomi launched and deployed perfectly, but a series of problems with the attitude control system caused the satellite to spin out of control and break up.

(920) The data from Hitomi, limited as it is, is enough to make astronomers re-think the role of plasma in galactic evolution, according to McNamara. "The plasma can be thought of forming an enormous atmosphere those envelopes (e) whole clusters of galaxies.

(921) These hot atmospheres represent (eb) the failure of the past—the failure of the universe to create (eb) bigger galaxies," he said.

(922) "But it's also the hope for the future. This is the raw material for the future growth of galaxies - which is (=) everything: stars, planets, and people

(923) . It's the raw material that in the next several billion years is going to make (eb) the next generation of suns and solar systems. And how rapidly that happens is governed by (e&eb) the black hole."

(924) The observations give researchers, for the first time, a direct measurement of (e) the turbulent speed of the hot plasma.

(925) "This measurement tells us how the enormous energy released by (e) supermassive black holes regulates (e&eb) the growth of the galaxy and the black hole itself," said McNamara. More information: The quiescent intracluster medium in the core of the Perseus cluster, Nature, DOI: 10.1038/nature18627 Journal reference: Nature search and more info website provided by: University of Waterloo search and more info website

**The quiescent intracluster medium in the core of the Perseus cluster Nature 535, 117–121 (07 July 2016) doi: 10.1038/nature18627 Published online 06 July 2016**

(926) Clusters of galaxies are (=) the most massive gravitationally bound objects in the Universe and are (=) still forming.

(927) They are thus important probes1 of (e) cosmological parameters and (e&eb) many astrophysical processes.

(928) However, knowledge of the dynamics of (e) the pervasive hot gas, the mass of which is much larger than *(e) the combined mass of all the stars in the cluster, is lacking.

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| (929) Such knowledge would enable (eb) insights into the injection of (e) mechanical energy by (e) the central supermassive black hole and the use of (e) hydrostatic equilibrium for (e) determining cluster masses. |
| (930) X-rays from the core of the Perseus cluster are emitted by (e) the 50-million-kelvin diffuse hot plasma filling (e) its gravitational potential well. |
| (931) The active galactic nucleus of the central galaxy NGC 1275 is pumping jetted energy into (e&eb) the surrounding intracluster medium, creating (eb) buoyant bubbles filled with (e) relativistic plasma. |
| (932) These bubbles probably induce *(eb) motions in the intracluster medium and heat the inner gas, preventing (e) runaway radiative cooling—a process known as (=) active galactic nucleus feedback2, 3, 4, 5, 6. |
| (933) Here authors report X-ray observations of (e) the core of the Perseus cluster, which reveal (eb) a remarkably quiescent atmosphere in which the gas has (e) a line-of-sight velocity dispersion of 164 ± 10 kilometres per second in the region 30–60 kiloparsecs from the central nucleus. |
| (934) A gradient in the line-of-sight velocity of 150 ± 70 kilometres per second is found across the 60-kiloparsec image of the cluster core. Turbulent pressure support in the gas is (=) four per cent of the thermodynamic pressure, with (e&eb) large-scale shear at most doubling this estimate. |
| (935) Authors infer that a total cluster mass determined from (e) hydrostatic equilibrium in (eb) a central region would require little correction for (e) turbulent pressure. |

**Dark matter guides growth of supermassive black holes February 18, 2015**

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| (936) This illustration (deleted attributable and ascribable to spatial restraints) shows two spiral galaxies - each with supermassive black holes at (eb) their center - as they are about to collide and form an elliptical galaxy. |
| (937) Every massive galaxy has (e) a black hole at its center, and the heftier the galaxy, (eb) the bigger its black hole. |
| (938) But why are the two related? After all, the black hole is millions of times smaller and less massive than (eb) its home galaxy. |
| (939) A new study of football-shaped collections of stars called elliptical galaxies provides (eb) new insights into the connection between a galaxy and (e&eb) its black hole. |
| (940) It finds that the invisible hand of dark matter somehow influences (e&eb) black hole growth. |
| (941) "There seems to be a mysterious link between the amount of dark matter a galaxy holds and (e&eb) the size of its central black hole, even though the two operate on (e&eb) vastly different scales," says lead author Akos Bogdan of the Harvard-Smithsonian Center for Astrophysics (CfA). |
| (942) This new research was designed to address a controversy in the field. Previous observations had found a relationship between the mass of the central black hole and (e&eb) the total mass of stars in elliptical galaxies. |
| (943) However, more recent studies have suggested (eb) a tight correlation between the masses of the black hole and (e&eb) the galaxy's dark matter halo. It wasn't clear which relationship dominated. |
| (944) In our universe, dark matter outweighs (e) normal matter - the everyday stuff we see all around us - by a factor of 6 to 1. |
| (945) We know dark matter exists only from (e) its gravitational effects. |
| (946) It holds together galaxies and (e&eb) galaxy clusters. |

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| (947) Every galaxy is surrounded by (e&eb) a halo of dark matter that weighs as (=) much as a trillion suns and extends for (e) hundreds of thousands of light-years. |
| (948) To investigate the link between dark matter halos and (e&eb) supermassive black holes, Bogdan and his colleague Andy Goulding (Princeton University) studied more than 3,000 elliptical galaxies. |
| (949) They used star motions as tracers to (e) weigh the galaxies' central black holes. |
| (950) X-ray measurements of hot gas surrounding (e&eb) the galaxies helped weigh the dark matter halo, because (e) the more dark matter a galaxy has, (eb) the more hot gas it can hold onto. |
| (951) They found a distinct relationship between the mass of the dark matter halo and (e&eb) the black hole mass - a relationship stronger than (e) that between a black hole and (e&eb) the galaxy's stars alone. |
| (952) This connection is likely to be related to (e&eb) how elliptical galaxies grow. |
| (953) An elliptical galaxy is formed when (e) smaller galaxies merge, their stars and dark matter mingling and (e&eb) mixing together. |
| (954) Because the dark matter outweighs (e) everything else, it molds (e&eb) the newly formed elliptical galaxy and guides (e&eb) the growth of the central black hole. |
| (955) "In effect, the act of merging creates a gravitational blueprint that the galaxy, the stars and the black hole will follow in order to build themselves," explains Bogdan.The paper describing this work has been accepted for publication in the Astrophysical Journal. This result relied on data from the Sloan Digital Sky Survey and the ROSAT X-ray satellite's all-sky survey. Journal reference: Astrophysical Journal search and more info website Provided by: Harvard-Smithsonian Center for Astrophysics search and more info website Read more at: http://phys.org/news/2015-02-dark-growth-supermassive-black-holes.html#jCp |

**Network theory sheds new light on origins of consciousness by Melanie Moran (APS site)**

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| (956) The black dots correspond to (e&eb) the 264 areas of the cerebral cortex that the researchers probed, and the lines correspond to (e&eb) the increased strength of the functional connections between each of these brain areas when subjects consciously perceive the … |
| (957) Where in (eb) your brain do you exist? Is your awareness of the world around you and of yourself as an individual the result of (e) specific, focused changes in (eb) your brain, or does that awareness come from (e) a broad network of neural activity? How does your brain produce (eb) awareness? |
| (958) Vanderbilt University researchers took a significant step toward answering these longstanding questions with a recent brain imaging study, in which they discovered global changes in how brain areas communicate with (e&eb) one another during awareness. Their findings, which were published March 9 in the Proceedings of the National Academy of Sciences, challenge previous theories that hypothesized much more restricted changes were responsible for (e) producing awareness. |
| (959) "Identifying the fingerprints of consciousness in humans would be (=) a significant advancement for basic and medical research, let alone its philosophical implications on (e&eb) the underpinnings of the human experience," said René Marois, professor and chair of psychology at Vanderbilt University and senior author of the study. "Many of the cognitive deficits observed in various neurological diseases may ultimately stem from (e) changes in how information is (=) communicated throughout the brain." |
| (960) Using graph theory, a branch of mathematics concerned with explaining the interactive links between (e&eb) members of a complex network, such as (=) social networks or flight routes, the researchers aimed to characterize (eb) how connections between the various parts of the brain were related to (e&eb) awareness. |
| (961) "With graph theory, one can ask questions about how efficiently the transportation networks in the United States and Europe are connected via (e&eb) transportation hubs like LaGuardia Airport in New |

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| York," Douglass Godwin, graduate student and lead author on the research, said. "We can ask those same questions about (e&eb) brain networks and hubs of neural communication." |
| (962) Modern theories of the neural basis of consciousness fall generally into two camps: focal and global. Focal theories contend there are specific areas of the brain that are critical for generating (eb) consciousness, while global theories argue consciousness arise from (e) large-scale brain changes in activity. This study applied graph theory analysis to (e&eb) adjudicates between these theories. |
| (963) The researchers recruited 24 members of the university community to participate in a functional magnetic resonance imaging (fMRI) experiment. While in the fMRI scanner, participants were asked to detect (e&eb) a disk that was briefly flashed on a screen. |
| (964) In each trial, participants responded whether they were able to detect (e&eb) the target disk and how much confidence they had in their answer. |
| (965) Experimenters then compared the results of the high-confidence trials during which the target was detected to the trials when (e) it was missed by participants. |
| (966) These were treated as (=) "aware" and "unaware" trials, respectively. |
| (967) Comparison of aware and unaware trials using (e) conventional fMRI analyses that assess the amplitude of brain activity showed (eb) a pattern of results typical of similar studies, with only a few areas of the brain showing (eb) more activity during detection of the target than when participants missed seeing it. |
| (968) The present study, however, was interested not simply in what regions might be more activated with (e&eb) awareness, but how they communicate with (e&eb) one another. |
| (969) Unlike the focal results seen using (e) more conventional analysis methods, the results via (e&eb) this network approach pointed toward (e&eb) a different conclusion. |
| (970) No one area or network of areas of the brain stood out as (=) particularly more connected during awareness of the target; the whole brain appeared to become (=) functionally more connected following reports of awareness. |
| (971) "We know there are numerous brain networks that control (e&eb) distinct cognitive functions such as (=) attention, language and control, with each node of a network densely interconnected with (e&eb) other nodes of the same network, but not (e) with other networks," Marois said. |
| (972) "Consciousness appears to break down (e) the modularity of these networks, as we observed (e) a broad increase in functional connectivity between these networks with awareness." |
| (973) The research suggests that consciousness is likely a product of (e) this widespread communication, and that we can only report things that we have seen once they are being represented in (eb) the brain in this manner. |
| (974) Thus, no one part of the brain is (=) truly the "seat of the soul," as René Descartes once wrote in a hypothesis about the pineal gland, but rather, consciousness appears to be (=) an emergent property of how information that needs to be acted upon (e&eb) gets propagated throughout the brain. |
| (975) "We take for granted how unified our experience of the world is. We don't (e) experience separate visual and auditory worlds, it's all integrated into (e&eb) a single conscious experience," Godwin said. |
| (976) "This widespread cross-network communication makes sense as (=) a mechanism by which consciousness gets integrated into (eb) that singular world." More information: Breakdown of the brain's functional network modularity with awareness, Douglass Godwin, DOI: 10.1073/pnas.1414466112 Journal reference: Proceedings of the National Academy of Sciences search and more info website: Provided by Vanderbilt University search and more info website Scientists find 'hidden brain signatures' of consciousness in vegetative state patients |

**Satellite studies x-ray emission from matter falling into the black hole at the center of a galaxy December 5, 2014 Figure 1: Photograph of galaxy NGC 3227. Credit: Ken Crawford (deleted due to spatial restrictions. Please excuse me. )**

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| (977) Most galaxies are assumed to have (e) at their heart a supermassive black hole that draws in (e) vast amounts of surrounding matter. |
| (978) As this matter is sucked in,(eb)it releases energy in the form of intense x-ray emissions that in some cases can be more intense than (e) the emission from all the stars in the galaxy combined. |
| (979) Through detailed study of the x-ray emissions from (e) the center of the galaxy known as NGC 3227, Hirofumi Noda from the RIKEN Nishina Center for (e) Accelerator-Based Science and colleagues have now revealed (eb) the multiple processes responsible for (e) these emissions. |
| (980) "Our results show the presence of multiple mechanisms of energy conversion of matter near (eb) a central black hole," explains Noda. |
| (981) The x-ray emissions from the centers of galaxies provide (eb) valuable information on the properties of (e) their central black holes, as well as (=) the surrounding matter and ultimately the history of galaxies in (eb) general. |
| (982) However, it can be difficult to separate (e&eb) the original radiation emitted by (e) matter close to (e) the black hole from (e) secondary processes caused by (e) other matter around (e&eb) the black hole. |
| (983) The typical spectral analysis applied to (e&eb) images of galactic centers is (=) unable to clearly distinguish the different emission processes. |
| (984) Noda and his colleagues instead studied the change in emissions over (e&eb) time. |
| (985) By analyzing the different radiation patterns and (e&eb) their correlations, they were able to separate primary emission from (e) the emission spectra arising from (e) secondary processes such as reflection. |
| (986) Using this approach and wide-range x-ray data obtained over (e&eb) the course of a few weeks by (e) the Japanese research satellite Suzaku, the team identified (e&eb) three distinct spectral components from (e) NGC 3227—a typical type I Seyfert galaxy that lies (eb) about 80 million light years from Earth (Fig. 1). |
| (987) The intensity of the source varied by (e) almost an order of magnitude over (e&eb) the course of the observations, and by analyzing these variations, the researchers could correlate the spectral components with (e&eb) three different mass accretion flows into (e&eb) the galaxy's black hole. |
| (988) Noda believes that studying the ultraviolet emissions from (e) matter at a greater distance from (e) the black hole could also provide (eb) further insight into the detailed physics of (e) supermassive black holes. X-ray emissions, however, are still expected to reveal (eb) much more information if (e) the spectral resolution of observations can be improved. |
| (989) "The soon-to-be-launched Japanese x-ray satellite ASTRO-H will achieve an unprecedentedly fine energy resolution, which will enable (eb) us to understand the properties of black holes, their evolution and their interaction with (e&eb) the galaxy around them." More information: Noda, H., Makishima, K., Yamada, S., Nakazawa, K., Sakurai, S. & Miyake, K. "Suzaku studies of the central engine in the typical type I Seyfert NGC 3227: Detection of multiple primary x-ray continua with distinct properties." Astrophysical Journal 794, 2 (2014). DOI: 10.1088/0004-637X/794/1/2Read more at: http://phys.org/news/2014-12-satellite-x-ray-emission-falling-black.html#jCp |

**Researchers chip away at Smale's 7th unsolved problem in mathematics July 15, 2016 by Lisa Zyga feature Thompson's problem**

(990) The disconnectivity graph for (e) 147 charges on (e&eb) a sphere has a structure-seeking “palm tree” organization.

(991) The five lowest minimum energy configurations are shown. Credit: Mehta et al. ©2016 American Physical Society

(992) (Phys.org)—How do you arrange a group of points on (e&eb) the surface of a sphere so that all the points are as far apart from (e) each other as possible?

(993) With two points, the answer is easy: place them on opposite sides of the sphere, as if they are (=)endpoints of the diameter. With three points, make them the vertices of an equilateral triangle, and so on. But as the number of points increases, so does the difficulty of the problem.

(994) This puzzle is the essence of Thomson's problem, which asks how to arrange equal charges (such as electrons) on (e&eb)the surface of a sphere in (e) a way that minimizes their electrostatic potential energy—the energy caused by (e) all of the electrons repelling each other.

(995) According to Coulomb's law, the total energy is inversely related to (e&eb) the sum of the distances between all possible pairs of charges, so (e) the goal is to spread the charges as far apart as (=) possible.

(996) This task is more difficult than it sounds—Thomson's problem has been rigorously solved only for (e) numbers of 2, 3, 4, 6, and 12 charges.

(997) In 1998, mathematician Steven Smale identified the problem of how to choose starting points close to (e) the lowest-energy state (which makes it easier to solve Thomson's problem) as (=) the seventh problem on his list of 18 unsolved problems for the 21st century.

(998) Part of the reason why Thomson's problem is so important is because (e) its applications are so far-reaching.

(999) In 1904, J.J. Thomson originally proposed the model of charges on a sphere to describe (eb) the structure of an atom.

(1000) Even though experiments disproved this "plum pudding model" long ago, the Thomson problem still has notable applications in (e&eb) chemistry (for understanding how electrons fill (e) electron shells in atoms), biology (for determining (eb) the arrangements of proteins on the shells of spherical viruses), as well as in physics, computer science, and such practical applications as determining (eb) the optimal placement of communication satellites around the Earth.

**Spheres on trees**

(1001) Now in a new paper published in Physical Review Letters, a team of mathematicians, engineers, and scientists from the US, the UK, and Australia has (e) taken a new approach to Thomson's problem that makes it much easier to determine (eb) the lowest-energy configuration.

(1002) They have constructed for seven numbers of charges (every third number from 132 to 150), tree-shaped disconnectivity graphs, where (e) the vertical axis or "trunk" corresponds to (e&eb) the energy of a particular charge arrangement.

(1003) Each "branch" terminates at (eb) a local minimum, which are (=) the states that have (e) lower energies than all of their neighboring states, and so they are (=) candidates for the ultimate lowest energy state, the global minimum.

(1004) By visualizing the problem in this way, the researchers noticed that these particular graphs don’t (e) have lots of branches extending from lots of (e) other branches.

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| (1005) Instead, every branch connects to (e&eb) only a few other branches and then to (e&eb) the trunk at regularly spaced energy thresholds, so that (e) the graph resembles a palm tree or single funnel structure. |
| (1006) The researchers found that this "funneled potential energy landscape" is (=) characteristic of a highly ordered structure and displays (eb) characteristics of a small-world network. |
| (1007) As a result, it provides (eb) an important clue for finding the global minimum. |
| (1008) It tells the researchers to start their optimization algorithms using (e) the local minima because, (e) in these networks, it turns out that every local minimum are (=) always within 5-7 steps (branches) of the global minimum. |
| (1009) This is true even for (e) local minima that have (e) much higher energies than the global minimum, and even when (e) the total number of local minima is very large. |

**Thompson's problem 2**

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| (1010) By visualizing the problem from a new perspective, the researchers found (eb) that lower-energy configurations have (e) more connections than (e) higher-energy configurations do. Credit: Mehta et al. ©2016 American Physical Society |
| (1011) The disconnectivity graph reveals (eb) other information, such as those lower-energy local minima have (e) more connections to (e&eb) other states than do higher-energy local minima. |
| **(1012)** The researchers also discovered that the global minimum is (=) always the most highly connected node in the entire network, making it the network's central node. |

**Implications for Thomson's and Smale's problems**

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| (1013) Using this insight from the network's single-funneled, small-world structure, finding the global minimum for Thomson's problem becomes (=) much easier than before. |
| (1014) "Our work looks at Smale's seventh problem from (e) a completely different perspective and (e&eb) sheds novel light on it," coauthor Dhagash Mehta, at the University of Notre Dame and the University of Adelaide, told Phys.org. "In this work, methods developed by the theoretical chemistry community have helped understand (eb) a deep mathematical problem. Often it is the other way around." |
| (1015) As the researchers explain, it's easier to solve Thomson's problem in these particular cases than it is to solve Smale's problem (of choosing good starting points). So although the results will likely be useful, they do not go very far toward solving Smale's seventh problem. |
| (1016) "I think 'chip away' is about right," said coauthor David Wales at University Chemical Laboratories in Cambridge, UK. "There is (=) no rigorous mathematical progress on the problem from (e) an analytic point of view." |
| (1017) In the future, the researchers plan to extend this approach to (e&eb) larger numbers of charges. |
| (1018) From earlier work, they expect that landscapes with (e&eb) more than 400 charges will start to display (eb) multiple funnels, so (e) the small-world structure may disappear. |
| (1019) "While we have only shown data for seven numbers, we have strong reasons to believe that the single funnel is (=) a feature for numbers less than 150," said coauthor Halim Kusumaatmaja at Durham University in Durham, UK. |
| (1020) "For larger numbers, there will likely be (=) multiple funnels. |
| (1021) Nonetheless, the network analysis could still be exploited to (e) help us quickly identify candidates for (e) the global minimum." |

(1022) Other lines of work include exploiting the small-world properties discovered here to improve other optimization algorithms and develop novel algorithms, as well as to incorporate weight and direction into these networks, which may provide additional insight into the Thomson problem.

(1023) "The social network analogy for networks of minima of (e) the Thomson problem will go further when (e) we analyze other network properties of these networks of minima," Mehta said.

(1024) "Our results will also help in constructing novel algorithms to find (eb) the global minimum more efficiently by (e) exploiting these network properties." More information: Dhagash Mehta et al. "Kinetic Transition Networks for the Thomson Problem and Smale's Seventh Problem." Physical Review Letters DOI: 10.1103/PhysRevLett.117.028301, Also at arXiv: 1605.08459 [cond-mat.soft] Journal reference: Physical Review Letters search and more info website © 2016 Phys.org

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_7t}{dt} = N_7(t-T_7)F_7\big(R_7(t-T_7)\big) - N_7(t)M_7\big(R_7(t)\big)$$

(9)

$$\frac{dR_7t}{dt} = B_7\big(R_7(t-T_7)\big) - D_7\big(R_7(t)\big) - N_7(t)G_7\big(R_7(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) and the effective prey regulation ($dD_7/dR_7 + NdG_7/dR_7 - dB_7/dR_7$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this

rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Look at the resemblance of String Theory and Abhinava's earlier theory and how the model could be applied: The world of manifestation and differentiated perceptions (Yikalpa) thus extends from one Centre to (eb) the next. **What is meant again is the realisation or the evolution of individual consciousness or general ledger to be (=)the cosmic general ledger is the sole criterion for the Satchidananda state where all dualities cease** Although it is never in fact divorced from (e) the subject who resides there, the ignorant fail (e) to grasp this fact and so, cut off from (e) the Centre, the world of objectivity becomes (=) for them the sphere of Maya. Bhagavatotpala quotes the Light of Consciousness amvitprakdsa): This ever pure experience (suddhanubhava) is variegated by (e) each form [revealed within it] Even so it remains unstained (nirmala) when moving to another. Just as a cloth which is naturally white, once dyed, cannot (e) change colour without [first] becoming white again, similarly the pure power of awareness, (citi) once coloured by (e) form, is pure [again] at the Centre where that form is (=) abandoned and from whence it proceeds to (e&eb) another. In his Essence of Vibration (Spandasarpdoha), K§emaraja explains that the rise and fall of every individual perception in the field of awareness is (=) a specific pulsation of consciousness. This is exactly the individual general ledger transaction lifted to the transaction in cosmic general ledger. During The Initial Instant Of Perception, T Consciousness Is Manifestly Apparent And The Yogi, Participating In Its Plenitude, Observes The Outer World Without Being Attached To Any Particular Divine Body And Sacred Circle Of The Senses Singling It Out From Any Other, Like A Man Who Observes A City From A High Mountain Peak.

He Sees The Outer World Reflected Within His Consciousness Free Of Thought-Constructs And So 'Stamps* The Outer On The Inner While Absorbing The Object And Means Of Knowledge In The Pure Subject Which Grasps Them As The Expansion Of His Own Nature. K§emaraja Says: By Penetrating Into Bhairavimudra, The Yogi Observes The Vast Totality Of Beings Rising From, And Dissolving Into, The Sky [Of Conscious- Ness], Like A Series Of Reflections Appearing And Disappearing Inside A Mirror. Through (e&eb) The Practice Of Bhairavimudra The Yogi Realises That He Is (=) The Substratum Consciousness (Adhisfhatr) Which Both Underlies (e&eb) And Is (=) The Essence Of All Things. He Discovers That Phenomena Have (e) No Independent Existence Apart From (e) Him And So Are, In This Sense, Void. At The Same Time, He Realises That Because (e) All Things Are Consciousness, They Are (=) Far From Unreal. He Views The Outer World Yet Sees (e) It Not. Beyond Both Voidness And Non-Voidness He Penetrates Into (e&eb) The Supreme Abode (Pararn Padam) Of Siva's Consciousness. [The Powers Of The Senses] Endowed With (e&eb) The Attributes Of The Great Union [Between Subject And (e&eb) Object] Whose Form Is The Wakening Of Man's Spiritual Potential (Kuntfalinf), Fill (e&eb) [With Consciousness] The Outer Clatter Of Diversity (Bhedatfambara) Born Of (e) Its Intense Power And Are Then Established In (eb) The Unobscured Abode Of The Void Of Consciousness To Shine [There] Eternally. Thus Residing Beyond (e) Being And Non-Being, The Sole Protector Of The Unity Which Is Tranquil And Expanding [Consciousness], Whose Glory Is All-Embracing And Form Unobscured, Is Called Bhairavimudra. Through (e&eb) The Practice Of Bhairavimudra, The Yogi Unites The Universal Vibration Of T Consciousness With (e&eb) The Individual Pulsation Of Objectivised 'This' Consciousness. The Two Aspects Of Consciousness Are Now In (eb) A State Of Equilibrium Like (=) The Two Pans Of An Evenly Weighed Balance And The Yogi Experiences (e&eb) The Pure Knowledge (Suddhavidya) That: 'I Alone Am (=)All Things'. Thus Becoming (=) The Master Of The Wheel Of Energies He Is (=) Free, Like Siva, To Create And Destroy. When [The Yogi] Is Well Established, Without (e) Wavering, Solely In The Integral Egoity Of His Authentic Nature, The Spanda Principle, And Is

Absorbed In Contemplation (Samavisfa), He Becomes (=) One With It The Doctrine Of Vibration (Tanmaya). Then . . . Dissolving And Creating The Universe By Means Of (er) His Introverted And Extroverted Absorption, He Destroys And Creates (e&eb) All Things Out Of Sarikara, His Innate Nature. [Thus] He Assumes The State Of The Universal Experiencer And Having Absorbed All That Is To Be Experienced From [The Grossest Level] — Earth — To (The Subtlest) — Siva — He Reaches (eb) The State Of The Supreme Subject By (e) Progressively Recognising [His Identity With Him]. Thus, Introverted And Extroverted Absorption Both Lead To (eb) The Recognition Of The Pulsation (Spanda) Of (e) One's Own Consciousness. This Is Where (e) Individual General Ledger Becomes (=) Cosmic General Ledger. At The Level Of Consciousness Corresponding To (e&eb) Siva's Basic State (Saryibhavdvastha), The Alternation From Inner To (e&eb) Outer Is Instantaneously Resolved Into (e&eb) The Vibration Of His Nature. When The Yogi Finally Comes To Be (=) Constantly Aware Of This Reality, His Enlightenment Is (=) Full And Perfect. Freed Of All Means (Anupaya) And Delighting In The Power Of His Bliss (Anandasakti), He Knows And Does Whatever He Pleases. The Yogi Seeking Self-Realisation Must Acquire (e) Mastery Over This Movement. Ksemaraja Stresses That The Doctrine Of Vibration Teaches That Liberation Can Only Be Achieved By (e) First Withdrawing All Sense Activity In Introverted Contemplation (Nimilanasamadhi) To (e) Then Experience The 'Great Expansion' (Mahdvikdsd) Of (e) Consciousness While (e) Recognising This To Be (=) A Spontaneous Process Within (eb) It. This Is Done Through (e&eb) The Practice Of Kramamudra. A Passage From The Now Lost Kramasiitra Explains: Although The Adept's Attention [May Be] Outwardly Directed, (eb) He Enjoys Contemplative Absorption Through (e&eb) The Introverted Aspect Of (e) Kramamudra. Initially He Turns Inward From (e)The Outside World And [Then] From Within [Himself] He Exits Into The Outer World Under The Influence Of His Absorption. There Is No Necessity Of Any Glorification Or Mortification Here. It Is Just Like Joining A Job And Retiring And Realising All That Happened Is Just By His Own Effort And Nothing Else. Thus The Sequence (Krama) In This Attitude (Mudra) [Ranges Through] Both Inner And Outer. The Yogi Must Pervade

The Surface Level Of Awareness (Yyutthana) With The Same Bliss He Experiences Plunging Into The Depths Of Con- Templative Absorption (Samadhi). Note Here That Samadhi Is A Death Like State And Still Being Matter And Information Stored You Experience Consciousness. Submerging Himself And Emerging Repeatedly From Samadhi, He Eventually Recognises That The Unity Of Consciousness Pervades Both States: The Best Of Yogi's, Who Has Achieved A State Of Complete Absorption Even When Risen From Meditation, [Inwardly] Vibrating Like A Drunkard In Blissful Inebriation From The After-Effects Of The Nectar Of Con- Templation, Sees All Things Dissolving In The Sky Of Consciousness Like A Cloud In The Autumn Sky. He Plunges Repeatedly within Himself Divine Body and Sacred Circle of the Senses and Becomes (=) Aware Of His Identity With (e&eb) Consciousness By (e) The Practice Of Introverted Contemplation. Thus Even When He Is Said To Have Risen From (e) Absorption, He Is (=) One With [His] Experience Of It. Ksemaraja Goes On To Explain That This Practice Is Called 'Mudra' Because (e) It Both Fills The Adept With (e&eb) Bliss (Mud) And Is Itself The Bliss Of (e) Consciousness. Moreover, It Dissolves (e) Away (Dra) All Bondage And 'Stamps' The Universe Of Experience With (e&eb)The Seal (Mudra) Of The Fourth State (Turiya) Of Enlightened Consciousness Beyond, And Including, The Three States Of Waking, Dreaming And Deep Sleep. It Is Called 'Krama' Because (e) It Is The Root Source Of (e) All Emanation And All Other Conscious Processes Which Succeed (e&eb) One Another In Ordered Sequence (Krama) And Is, (=)At The Same Time, Their Successive (Krama) Appearance As Well. By The Practice Of Kramamudra The Opposites Fuse And Siva And Sakti Unite. They Yogi Comes To Experience The Simultaneous Pervasion Of All The Lower, Grosser Categories Of Existence By The Higher And The Presence Of The Lower In The Higher. Commencing His Practice In A Low Form Of Bhairavimudra, The Yogi Conjoins The Outer With The Inner; Then, In Kramamudra, He Fills Both The Outer With (e&eb) The Inner And (e&eb) The Inner With (e&eb)The Outer. When He Achieves (eb) Perfection In This Two-Fold Movement, He Attains To (e) The Highest Form Of Bhairavimudra In Which The Two Merge Completely In (eb)The Experience Of (e) The Absolute (Anuttara), Free

Of (e) All Differentiation And (e&eb) Polarities. If He Fails To Maintain Awareness Of This State, He Again Falls Into (eb) Kramamudra Until He Has Finally Completely Merged All The Highest States In (e&eb) The Lower And The Lower In (e&eb) The Higher. He Then No Longer Needs To (e) Resort To (e) Any Means (Anupaya) To (e) Achieve Liberation. All He Says or Does Anything He Perceives or Thinks, Instantly Occasions In (eb) Him The Highest Level Of Consciousness. Thus The Fruit Of Bhairavimudra Is (*=)The Wonder (Camatkara) Or Amazement (Vismaya) That Overcomes (e)The Yogi When (e) He Reaches The Plane Of (e) Union (Yogabhumika) Y Xsl Where (e) All Opposites Merge In (e&eb)The Radiance Of The Great Light Of Consciousness. The Stanzas On Vibration Teach: How Can One Who, As If Astonished, Beholds (e) His Own Nature As (=) That Which Sustains [The Existence Of Everything] Be Subject To (e&eb) This Painful Round Of Transmigration (Kusrti)? The Yogi, Recognising His True Nature To Be (=)The Supreme Subject, Is Astonished To (e) Suddenly Discover That The Individual He Thought He Was, Caught Up (e&eb) In The Trammels Of Thought And Living In A World Enmeshed In (e&eb) The Web Of Time And Space, Does Not Really Exist At All. He Experiences A 'Turning About' (Paravrtti) In (eb)The Deepest Seat Of The Doctrine Of Vibration Consciousness As He Penetrates His True Nature. The Sudden Eruption Of This Intuition (Pratibha) Arouses In Him A Cry Of Amazement As He Transcends All Thought-Constructs And, Perfectly Absorbed In His Own Nature, Is Liberated. The Path To Liberation Essentially, Spanda Doctrine Is Concerned With Two Matters. The First Is To Impart To Those Who Are Fit To Receive The Teachings A Deeper Understanding Of The Ultimate Goal Of Life (Upeya). When We Have Understood What Truly Benefits Us And Is Worth Attaining And What, On The Contrary, Is Of No Real Value But Stands In The Way Of This Attainment, We Can Begin To Make Progress Towards Our Goal. This Is Spanda Doctrine's Second Concern, Namely, To Show The Way In Which We Can Develop Spiritually Through Siva's Grace And The Right Application Of The Means To Realisation That It Teaches. When Both These Aspects Of The Teaching Have Been Correctly Understood And Applied, The Spanda Yogi Achieves A Clear And Permanent

Realisation Of His Goal And Is Liberated, Thus Fulfilling The Ultimate Aim Of The Teaching. The Doctrine Of Vibration Is Not Meant For The Spiritually Dull. It Is Not For The Worldly Whose Consciousness, Clouded By Ignorance, Is As If Dreaming, Even During The Waking State Of Daily Life, The Dream Of Its Own Thought-Constructs. The Teachings Are Meant For Those Who Are Awake (Prabuddha), Those Who, Full Of Faith And Reverence, Are Always Alert And Intent On Discerning The True Nature Of Ultimate Reality. This Reality Is Understood In Three Basic Ways. The First Is Purely Transcendental. The Stanzas Choose This Aspect As The One Which Formally Defines It Most Specifically. Ultimate Reality Transcends All The Opposites, Including Subject And Object. This Does Not Mean, However, That It Is An Unconscious Void, A Mere Absence Of All Existence. In Fact, This Negative Characterization Of Reality (Which Includes Also A Denial Of All That Is Unconscious) Implies A Positive Immanence In Which The Opposites Are United In The Oneness Of Pure Consciousness That Is Equally Siva And Spanda, His Universal Activity. These Two Seemingly Contrasting Aspects Are Reconciled In The Third, Namely, Reality Understood As The Essential Nature Of All Things. Although Universal And Everywhere The Same, It Is Understood To Be The Essential And Specific Nature Of Each Existent As Its 'Own Nature' (Svabhava).

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_7/dR_7$ is a measure of the control acting through the debit delay and setting $dB_7/dR_7 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_7/dR_7$ is a measure of the control acting through the credit delay although setting $dF_7/dR_7 = 0$ does not negate the effect of a predator delay; the value of $M_7$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_7/dR_7$and normalized by $M_7$ (see7a). The absolute stability of such

a system is summarized in Fig. 7 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

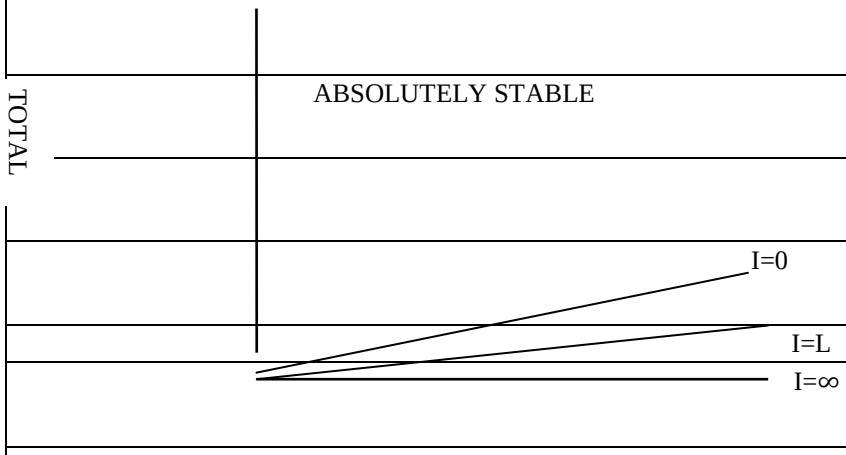
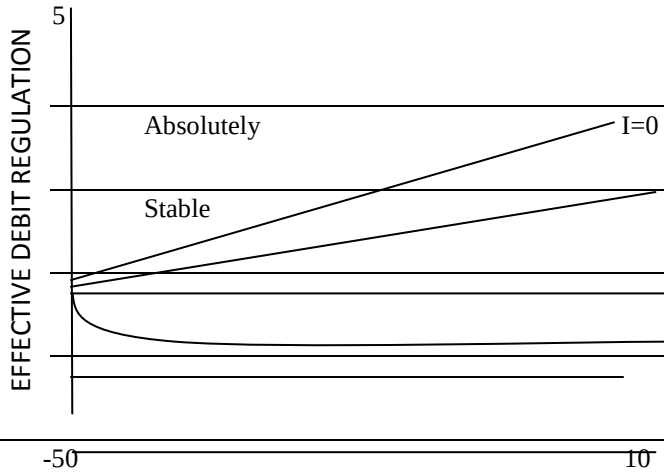



FIG7. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_7/dR_7 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




**DELAYED CREDIT EFFECT (LACK OF BALANCE)**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_7/dR_7 + N_7 dG_7/dR_7 + dB_7/dR_7)/M_7$). The immediate ($I_7$) and lagged ($L_7$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_7$(see 7b), is increased relative to the lagged part, $L_7$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_7 > I_7$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_7$is somewhat large than $I_7$ then the condition constraining prey regulation of recruitment ($dB_7/dR_7 < 0$) is

$$\frac{dD_7}{dR_7} + N_7\frac{dG_7}{dR_7} + \frac{dB_7}{dR_7} > \frac{(L_7 - I_7)M_7}{2} \qquad (10)$$

If $L_7 < I_7$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

This equation is of paramount importance in the following cases: According to the principles of geometric quantum mechanics, the physical characteristics of a given quantum system can be represented by geometrical features that are preferentially identified in this complex manifold. Here we construct a number of examples of such features as they arise in the state spaces for spin, spin 1, spin and spin 2 systems, and for pairs of spin systems. A study is then undertaken on the geometry of entangled states. A locally invariant measure is assigned to the degree of entanglement of a given state for a general multi-particle system, and the properties of this measure are analysed for the entangled states of a pair of spin particles. With the specification of a quantum Hamiltonian, the resulting Schrödinger trajectories induce an isometry of the Fubini–Study manifold, and hence also an isometry of each of the energy surfaces generated by level values of the expectation of the Hamiltonian. For a generic quantum evolution,

the corresponding Killing trajectory is quasiergodic on a toroidal subspace of the energy surface through the initial state. When a dynamical trajectory is lifted orthogonally to Hilbert space, it induces a geometric phase shift on the wave function. The uncertainty of an observable in a given state is the length of the gradient vector of the level surface of the expectation of the observable in that state; a fact that allows us to calculate higher orders corrections to the Heisenberg relations. A general mixed state is determined by a probability density function on the state space, for which the associated first moment is the density matrix. The advantage of a general state is in its applicability in various attempts to go beyond the standard quantum theory, some of which admit a natural phase-space characterization. **Journal of Geometry and Physics Volume 38, Issue 1, April 2001, Pages 19–53 Geometric quantum mechanics Dorje C. Brody Lane P. Hughston DOI: 10.1016/S0393-0440(00)00052-8** Expansion (unmesa) of [consciousness] or the creative intuition {pratibha) is (=) [experienced] in the interval which divides two [moments] of differentiated perception (vikalpa) It is here that they arise and disappear. The Sastras and Agamas proclaim with reasoned argument that it is (=) free of thought-constructs {Nirvikalpa) and precedes (e) Path to Liberation all mental representations of any object. None can deny that a gap exists between perceptions insofar as two moments of thought are (=) invariably divided. This [gap] is (=) the undifferentiated unity of all the countless manifestations. Similarly, in the outer more objective sphere, where change consists (e) of the alterations in the configurations of manifest appearances (Abhasa\ the transition from one to (e&eb) another corresponds to (e&eb) a phase of pure luminosity that marks (eb) the beginning of one form and the end of another

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_7$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_7$

The - $D_7$ positions are defined by letting $\lambda_7 = iZ_7$ and separating (A1.2) into real and imaginary parts:

$$Z_7^2 + V_7 Z_7 \int_0^{\infty} \sin(Z_7 T_7) f(T_7) dT_7 - (X_7 + Y_7) = 0$$

$$Z_7 \left( U_7 - V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7 \right) = 0$$

Noting that $X_7 + Y_7 = 0$ is we assume X+Y=0.It follows that $Z_7 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_7 = V_7 \int_0^{\infty} cos\left(Z_7 T_7\right) f\left(T_7\right) dT_7$$

If $U_7 > |V_7|$then given the constraints of $f(T_7)$, $(A1.5) can\ never\ be\ satisfied\ Co$nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A7.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_7^2 + (U_7 - V_7 + M_7)\lambda_7 - M_7 \lambda_7 e^{-T_7 \lambda_7} + [(U_7 - V_7) M_7 + Y_7] - [((U_7 -$$

V 7M 7−]e−T 7λ 7=0

(For definitions see Appendix7) We follow the technique outlined in Appendix 7 and let $\lambda_7 = iZ_7$ and separate (A2.7) into real and imaginary parts

$$\boldsymbol{Z}_7^2 + M_7 Z_7 \sin(Z_7 T_7) - (A_7 M_7 + Y_7) + (A_7 M_7 + X_7) \cos(Z_7 T_7) = 0$$

$$(A_7 M_7) Z_7 - M_7 Z_7 \cos(Z_7 T_7) + (A_7 M_7 - X_7) \sin(Z_7 T_7) = 0$$

The warrant of Truth is ever elusive when we deconstruct the foundational justifications of our convenient notions about the way the world works. It is easy to confuse what is actually the creation of beliefs with the "discovery of Truth," a common goal of science and theology. Science offers no ontological Scientific Picture of the World but substitutes a number of applied theories as meta-theory. It's close, powerful and useful: it 'works' and allows use to do work. But the still mysterious "true" metatheory may reveal the basis of consciousness. But who knows when we will lift that veil? Permeating the living reality of our culture are certain contagious notions, fads and trends that have the ability to influence the way we think about the nature of Reality and ourselves. Some of them are toxic and can consume you; others are just whims of pop culture infectious and benign as the common cold. **Jung** described the concept of psychic contagion by certain archetypal forces inherent in the human psyche, which manifest in our spiritual lives and belief systems. He spoke of both conscious and unconscious contamination. Notions like this range from simple superstitions to scientific concepts, to urban myths. Notions sweep through our culture and insinuate themselves within its fabric, as fads, whether they are "real" or not, they can be influential. They mood alter us, make us feel we belong. An analysis of these notions is useful in distinguishing a common human phenomenon from any potential "alien influence" which may or may not be exerted on us from an unknown source. **Strange Attractors: Transference, Holography, and an Archetype Burke, J. (2003).**

**Strange Attractors: Transference, Holography, and an Archetype (Doctoral dissertation, Pacifica Graduate Institute, 2003).** In consideration to the Nobel Proze for Chemistry this year(2014), a stability nalysis of the proble in to single molexues would be a contemporaneous question that has to be taken cognisace of and given importance to . Determination of molecular structure by geometry optimization became routine only after efficient methods for calculating the first derivatives of the energy with respect to all atomic coordinates became available. Evaluation of the related second derivatives allows the prediction of vibrational frequencies if harmonic motion is estimated. More importantly, it allows for the characterization of stationary points. The frequencies are related to the eigenvalues of the Hessian matrix, which contains second derivatives. If the eigenvalues are all positive, then the frequencies are all real and the stationary point is a local minimum. If one eigenvalue is negative (i.e., an imaginary frequency), then the stationary point is a transition structure. If more than one eigenvalue is negative, then the stationary point is a more complex one, and is usually of little interest. When one of these is found, it is necessary to move the search away from it if the experimenter is looking solely for local minima and transition structures. The total energy is determined by approximate solutions of the time-dependent Schrödinger equation, usually with no relativistic terms included, and by making use of the Born–Oppenheimer approximation, which allows for the separation of electronic and nuclear motions, thereby simplifying the Schrödinger equation. This leads to the evaluation of the total energy as a sum of the electronic energy at fixed nuclei positions and the repulsion energy of the nuclei. Notable exceptions are certain approaches called direct quantum chemistry, which treat electrons and nuclei on a common footing. Density functional methods and semi-empirical methods are variants on the major theme. **Orlando Alvarez** discusses Polyakov's quantization of the string in the presence of a boundary allowing for an arbitrary topology for the world sheet. In addition to the dynamical conformal factor discovered by Polyakov, there are a finite

number of new degrees of freedom if the surface is more complicated than a sphere or a disc. The quantization of the Liouville theory in an arbitrary topology is also discussed. A one-loop calculation shows that the model is renormalizable if one performs a mass renormalization and an additive field renormalization. The renormalization group equations have a perturbative infrared unstable fixed point in all topologies. Copyright © 1983 Published by Elsevier B.V. **Nuclear Physics B Volume 216, Issue 1, 25 April 1983, Pages 125–184 Theory of strings with boundaries: Fluctuations, topology and quantum geometry Orlando Alvarez DOI: 10.1016/0550-3213(83)90490-X** Quantum geometry (the modern loop quantum gravity involving graphs and spin-networks instead of the loops) provides microscopic degrees of freedom that account for black-hole entropy. However, the procedure for state counting used in the literature contains an error and the number of the relevant horizon states is underestimated. In our paper a correct method of counting is presented by **Marcin Domagala and Jerzy Lewandowski.** Results lead to a revision of the literature of the subject. It turns out that the contribution of spins greater than 1/2 to the entropy is not negligible. Hence, the value of the Barbero–Immirzi parameter involved in the spectra of all the geometric and physical operators in this theory is different than previously derived. Also, the conjectured relation between quantum geometry and the black-hole quasi-normal modes should be understood again. **Marcin Domagala and Jerzy Lewandowski 2004 Class. Quantum Grav 21 5233 doi:10.1088/0264-9381/21/22/014 Black-hole entropy from quantum geometry** Quantum geometry predicts that a universe evolves through an inflationary phase at small volume before exiting gracefully into a standard Friedmann phase. This does not require the introduction of additional matter fields with ad hoc potentials; rather, it occurs because of a quantum gravity modification of the kinetic part of ordinary matter Hamiltonians. Authors draw the cognizance of application of the same mechanism that can explain why the present day cosmological acceleration is so tiny. DOI: http://dx.doi.org/10.1103/PhysRevLett.89.261301

**Inflation from Quantum Geometry Phys. Rev. Lett 89, 261301 – Published 12 December 2002 Martin Bojowald** The loop quantum cosmology of the closed isotropic model is studied with special emphasis on a comparison with traditional results obtained in the Wheeler-DeWitt approach by **Martin Bojowald** .Investigation includes the relation of the dynamical initial conditions to boundary conditions such as the no-boundary or the tunneling proposal and a discussion of inflation from quantum cosmology. **DOI: http://dx.doi.org/10.1103/PhysRevD.67.124023 Loop quantum cosmology, boundary proposals, and inflation Phys. Rev. D 67, 124023 – Published 23 June 2003 Martin Bojowald and Kevin Vandersloot** In loop quantum cosmology, the universe avoids a big bang singularity and undergoes an early and short super-inflation phase. During super-inflation, non-perturbative quantum corrections to the dynamics drive an inflaton field up its potential hill, thus setting the initial conditions for standard inflation. **Shinji Tsujikawa et al** show that this effect can raise the inflaton high enough to achieve sufficient e-foldings in the standard inflation era. They also analyse the cosmological perturbations generated when slow-roll is **violated after super-inflation** and show that loop quantum effects can in principle leave an **indirect signature** on the largest scales in the CMB, with some loss of power and running of the spectral index. **Shinji Tsujikawa et al 2004 Class. Quantum Grav 21 5767 doi:10.1088/0264-9381/21/24/006 Loop quantum gravity effects on inflation and the CMB Quantum geometry from quantum information** is the study of **Florian Girelli and Etera R Livine .**Loop quantum gravity defines the quantum states of space geometry as spin networks and describes their evolution in time. We reformulate spin networks in terms of harmonic oscillators and show how the holographic degrees of freedom of the theory are described as matrix models. This allows us to make a link with non-commutative geometry and to look at the issue of the semi-classical limit of loop quantum gravity from a new perspective. This work is thought of as part of a bigger project of describing quantum geometry in quantum information terms.

**Florian Girelli and Etera R Livine 2005 Class. Quantum Grav 22 3295 doi:10.1088/0264-9381/22/16/011 Reconstructing quantum geometry from quantum information: spin networks as harmonic oscillators.** Quantum geometric mechanics is the topic of study by **Dorje C. Brody Lane P. Hughston** establishing a clear relation between the quantum mechanics and quantum geometry and also the entangled states. The manifold of pure quantum states can be regarded as a complex projective space endowed with the unitary-invariant Fubini–Study metric. Note how differential gheometry concepts with Nunney's model could be entangle so that it is feasibly and viably applied to various type of model including consciousness (See Doctrine of Vibration)

Where $A_7 = U_7 - V_7$Using (A2.3), substitute for $\boldsymbol{Z}_7^2$ and $Z_7$in (A2.2):

$W_7^2 (X_7^2$-3 $M_7\ A_7\ X_7$- $M_7^2 (X_7$- $Y_7) + 2\ M_7^3\ A_7$+2 $A_7^2\ M_7^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_7$=1-cos ($Z_7 T_7$).

In the absence of a delay the system is stable for $A_7 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_7 > (X_7 - Y_7)/2\ M_7$ Thus a necessary condition for absolute stability is

$A_7 > [\ X_7 - Y_7/2\ M_7$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_7$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_7$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_7, A_7) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_7 \le 2$ *is* of interest First consider the extreme values

$$Q_7(0, A_7) = A_7^2\ (\mathrm{X}_7 - Y_7)$$

$$Q_7(2, A_7) = \left(A_7 + 2\,\mathrm{M}_7\right)^2 (2A_7\,\mathrm{M}_7 - (\mathrm{X}_7 - Y_7))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_7}{\partial A_7} = W_7^2\left(2M_7^3 + 4A_7 M_7^2 - 3M_7 X_7\right) + W_7\left(3M_7 A_7^2 - 2A_7 X_7 + 2M_7\left(2X_7 + Y_7\right)\right)$$

$$+2A_7\left(X_7 + Y_7\right)$$

$$=2\,M_7^2 W_7^2 (M_7 + \mathrm{A}_7) + M_7\,A_7^2 W_7 + \left(2A_7 M_7 - X_7 + Y_7\right)(A_7 W_7 + M_7 W_7^2$$

$$+ (2- W_7)((2\,M_7\,X_7 + M_7\,Y_7)\,W_7 + A_7\,(X_7 + Y_7)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_7$quadratic in $\mathrm{W}_7$touches the $W$ –axis only at its minimum value in which case (A2.9) shows that

further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_7$ has no minimum in the range $0 \le W_7 \le 2$. Defining

$Q_7(W_7, A_7) = a_7 W_7^2 + c_7$ *where* $a_7\ b_7$ *and* $c_7$ Are functions of $A_7$ we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_7 \ge b_7 \ge 0$$

And

$$b_7^2 \le 4a_7 c_7$$

$$Case(i): Y_7 = 0$$

*Let* $Y_7 = 0$ And try the substitution A= $X_7 = /2M_7\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_7 = 1$ The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}(W_7 X_7 / 2M_7)$ is positive for *W<2. Since* $Q_7$ increases with $A_7$ (see A2.9), increasing $A_7$ will generate no further solutions of $Q_7$=0. Hence the condition A>X/2M defines absolute stability when $Y_7 = 0$.

$Case\ (ii): Y_7 = y_7 X_7$ ,

Let $Y_7 = y_7 X_7$ And try the substitution $A_7 = X_7/(2 + A_7 y_7)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_7/(2 + n_7 y_7)$ we have

$$W_7^2 y_7 \left[X_7 n_7 \frac{(1 + n_7 y_7)}{(2 + n_7 y_7)} + M_7^2 (2 - n_7 + n_7 y_7)\right]$$

$$- W_7 X_7 \left[\frac{X_7 (1 + n_7 y_7)}{M_7^2 (2 + n_7 y_7)^2} + 2 y_7 (n_7 - 1) + \frac{X_7^2 (7 + y_7)}{M_7^2 (2 + n_7 y_7)}\right] = 0. (A2.12)$$

Since $n_7$ and $A_7$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_7$

Defining $E_7 = (2 + n_7 y_7)$, *the condition*(A2.77?)*can be written as*

$[\frac{X_7(E_7-7)}{E_7^2}]^2[\frac{X_7^2}{M_7^2} - 4 y_7 E_7^2 \frac{X_7}{M_7^2} + 4y_7 (n_7 - 2)E_7^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_7 / W_7^2)$ can be expressed as

$2 E_7 [y_7 E_7 + (y_7 E_7 (7 + y_7)(E_7 - n_7))^{7/2}] > \frac{X_7}{M_7^2} > 2E_7 [y_7 E_7 - y_7 E_7 (7 + y_7)(E_7 - n_7))^{7/2}]$

**ABSOLUTE STABILITY**

Hence if $X_7 / M_7^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0 $W_7$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_7-7)}{E_7^2}\left(\frac{X_7}{M_7^2}\right)^2-2y_7\frac{(E_7+n_7^2y_7)}{E_7}\left(\frac{X_7}{M_7^2}\right)-4y_7(E_7-n_7)$$

And has the solution

$$\frac{X_7}{M_7^2}>\frac{E_7}{(E_7-7)}[y_7(E_7+n_7^2y_7)+[y_7^2(E_7+n_7^2y_7)^2+4y_7(E_7-7)(E_7-n_7)]^{7/2}]$$

Given $E_7>n_7$this lower limit to $X_7/M_7^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_7/M_7^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_7$in the appropriate range and this limit is further lowered by (A2.74)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_7<n_7$ $(ie.,2A_7M_7^2<X_7-Y_7$)then (A2.7)and (A2.8)show that $Q_7$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_7}{M_7^2}>2E_7[y_7E_7-(y_7E_7(7+y_7)(E_7-n_7)]^{7/2}]$$

And

$$E_7>n_7$$

Replacing $E_7$ and $n_7$ in these conditions leads to the form

If $A_7^2 > 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.

If $A_7^2 < 2Y_7$ then $A_7 > (X_7 - Y_7)/2M_7$.+ $(Y_7 - A_7^2/2)^2/2M_7(X_7 - Y_7)$.

Note that by normalizing $A_7$ by $M_7$ and $X_7$ and $Y_7$ by $W_7^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_7^2 + (M_7 + U_7)\lambda_7 - (M_7 + V_7)\lambda_7 e^{2T_7\lambda_7} + (M_7 U_7 + Y_7)$$

$$-[M_7(U_7 + V_7) - X_7]\, e^{2T_7\lambda_7} + M_7 V_7\, e^{2T_7\lambda_7} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_7 = \mathrm{i} Z_7$ *and* separate

(A3.1)into real and imaginary parts

$$Z_7^2 + (M_7 + V_7) Z_7 \sin(Z_7 T_7) - (U_7 M_7 + Y_7) + (M_7(U_7 + V_7) - X_7)$$

$$\mathrm{X} \cos(Z_7 T_7) - 2M_7 V_7 \cos^2(Z_7 T_7) + M_7 V_7 = 0$$

$$(M_7 + U_7) Z_7 - (M_7 + V_7) Z_7 \cos(Z_7 T_7) + (M_7(U_7 + V_7) - X)$$

$$\mathrm{X} \sin((Z_7 T_7) - M_7 V_7 \sin(Z_7 T_7) = 0$$

Using (A.3.3) to eliminate $Z_7$and $Z_7^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_7 = 1 - cos\, Z_7\ T_7$and $A_7 = U_7 - V_7$we obtain

$$4\,M_7^2 V_7^2 W_7^4 + 2\,M_7 V_7\ W_7^3[(M_7 - V_7)^2 + A_7(M_7 + V_7) + 2(A_7 M_7 - X_7)]$$

$$+ W_7^2[M_7 - V_7)^2(2\,A_7\ M_7 + Y_7 - X_7) + (2A_7\ M_7 - X_7)^2 - A_7(M_7 - V_7)(2A_7\ M_7 - X_7)$$

$$+ M_7 V_7\,(A_7^2 + 4(X_7 + Y_7))]$$

$$+ W_7[A_7^2 + 2X_7)(2\,A_7\ M_7 - X_7) + 2A_7 Y_7(M_7 + V_7) - M_7 A_7^3] + A_7^2(X_7\ M_7 - Y_7) = 0$$

Being a quartic in $W_7$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_7 > 0$$

$$and$$

$$C_7 > (X_7 - Y_7)/2\,M_7$$

$recalling\ that\ A_7 = U_7$and defining $C_7 = U_7 + V_7$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_7(0, U_7 V_7 = A_7^2(X_7 + Y_7)$$

$$Q_7\left(2, U_7 V_7\right) = \left(C_7 + 2M_7^2\right)^2 (2C_7 M_7 + Y_7 - X_7)$$

So that given (A3.6) the extreme $W_7$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_7$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_7$the existence of a D-position at this point is determined by letting cos $(Z_7 T_7) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 e^{-\tau_7\lambda_7} + X_7 + Y_7 = 0 \qquad (A1.1)$$

Where $V_7 = dB_7/dR_7$ $U_7 = dD_7/dR_7 + N_7 dG_7/dR_7$ $X_7 = N_7 G_7 dF_7/dR_7$ and $Y_7 = -N_7 G_7 dM_7/dR_7$ (evaluated at the equilibrium). Note that $U_7 X_7$ and $Y_7$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_7$ as $(T_7)$ . (It follows that $f(T_7)$ 0 for all $T$, and $\int_0^\infty f(T_7)dt = 7.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_7^2 + U_7\lambda_7 - V_7\lambda_7 \int_0^\infty e^{-\tau_7\lambda_7} f(T_7) + X_7 + Y_7 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

## Absolute Stability in Predator- Prey Models

## APPLIED TO ACCETUATION ATTRITION MODELS

**Module numbered five**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems. The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itself vary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This is the condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference

between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \quad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t - T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\left(R(t)\right)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

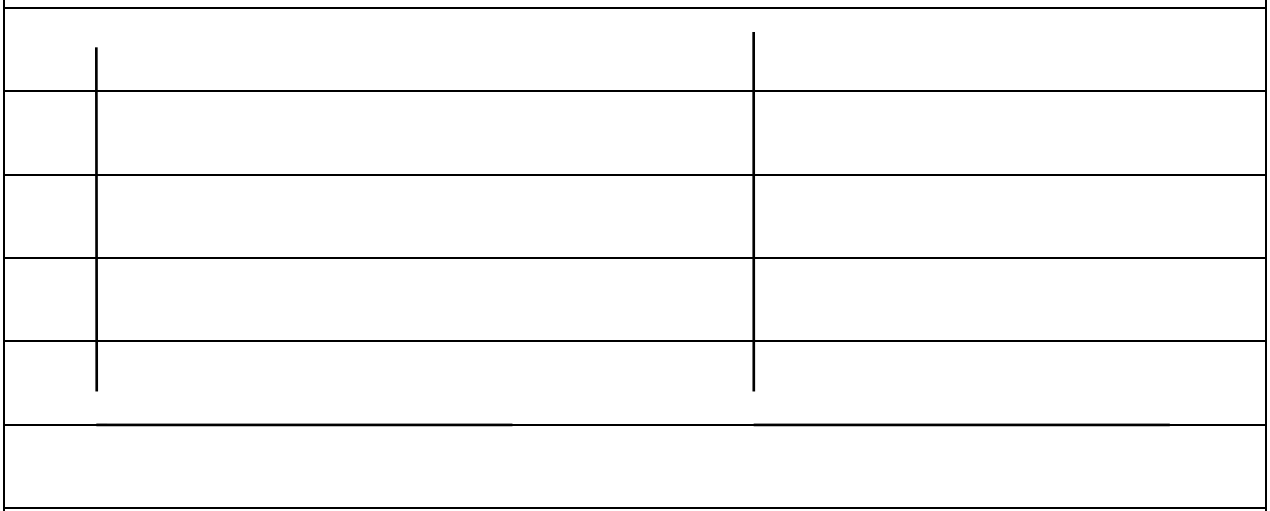
Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time

delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the



TIME DELAY

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period,

which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

As it is usually expressed outside of phenomenological texts, to undergo a conscious experience necessarily means that there is something it is like for the subject to have that experience (Nagel 1974; Searle 1992). This is obviously true of bodily sensations like pain. But it is also the case for perceptual experiences, experiences of desiring, feeling, and thinking. There is something it is like to taste chocolate, and this is different from what it is like to remember what it is like to taste chocolate, or to smell vanilla, to run, to stand still, to feel envious, nervous, depressed or happy, or to entertain an abstract belief. Yet, at the same time, as I live through these differences, there is something experiential that is, in some sense, the same, namely, their distinct first-personal character. All the experiences are characterized by a quality of mineness or for-me-ness, the fact that it is I who am having these experiences. All the experiences are given (at least tacitly) as my experiences, as experiences I am undergoing or living through. All of

this suggests that first-person experience presents me with an immediate and non-observational access to myself, and that consequently (phenomenal) consciousness consequently entails a (minimal) form of self-consciousness. To put it differently, unless a mental process is pre-reflectively self-conscious there will be nothing it is like to undergo the process, and it therefore cannot be a phenomenally conscious process. The mineness in question is not a quality like being scarlet, sour or soft. It doesn't refer to a specific experiential content, to a specific what; nor does it refer to the diachronic or synchronic sum of such content, or to some other relation that might obtain between the contents in question. Rather, it refers to the distinct givenness or the how it feels of experience. It refers to the first-personal presence or character of experience. It refers to the fact that the experiences I am living through are given differently (but not necessarily better) to me than to anybody else. It could consequently be claimed that anybody who denies the for-me-ness of experience simply fails to recognize an essential constitutive aspect of experience. Such a denial would be tantamount to a denial of the first-person perspective. It would entail the view that my own mind is either not given to me at all — I would be mind- or self-blind — or is presented to me in exactly the same way as the minds of others. There are also lines of argumentation in contemporary analytical philosophy of mind that is close to and consistent with the phenomenological conception of pre-reflective self-awareness. Alvin Goldman provides an example: [Consider] the case of thinking about x or attending to x. In the process of thinking about x there is already an implicit awareness that one is thinking about x. There is no need for reflection here, for taking a step back from thinking about x in order to examine it…When we are thinking about x, the mind is focused on x, not on our thinking of x. Nevertheless, the process of thinking about x carries with it a non-reflective self-awareness (Goldman 1970, 96). A similar view has been defended by Owen Flanagan, who not only argues that consciousness involves self-consciousness in the weak sense that there is something it is like for the subject to have the experience, but also speaks of the low-level self-consciousness involved in experiencing my experiences as mine (Flanagan 1992, 194). As Flanagan quite correctly points out, this primary type of self-consciousness should not be confused with the

much stronger notion of self-consciousness that is in play when we are thinking about our own narrative self. The latter form of reflective self-consciousness presupposes both conceptual knowledge and narrative competence. It requires maturation and socialization, and the ability to access and issue reports about the states, traits, dispositions that make one the person one is. To claim that every kind of self-consciousness is conceptual is overly cognitive. Bermúdez (1998), to mention one further philosopher in the analytic tradition, argues that there are a variety of **nonconceptual forms of self-consciousness** that are "logically and ontogenetically more primitive than the higher forms of self-consciousness that are usually the focus of philosophical debate" (1998, 274; also see Poellner 2003). This growing consensus across philosophical studies supports the phenomenological view of pre-reflective self-consciousness. (References: Stanford Encyclopedia, Post Modersnist Theories, Internet Encyclopedia of Philosophy) Korzybski pointed out that there are two ways to slice easily through life; to believe everything or to doubt everything. Both ways save us from thinking. So, from time to time you have to give yourself a Virtual Reality Check. The hallmark of postmodern philosophy has been disbelief or skepticism of all "metanarratives," or translations of reality. Postmodernism has even turned its profound skepticism on such important humanist concepts as "objective truth" and reason.Yet, for a deconstructionist postmodern society, individually we are still riddled with superstition and gullibility, and open to manipulation through our belief systems as any politician, philosopher, clergy, or salesperson will attest. Further, most people are painfully naive when it comes to even the simplest scientific understanding. Most of us don't have a clue about the fundamental nature of physical reality or our own psychological nature, and our ability to be fooled by our senses and mind, or methods of social and technological persuasion. We may think we know, but the penetration is deeper than we dare imagine. Nature rejects the naiveté that seeks absolute truth. We are beginning to realize, individually and culturally, that "realities" are all human constructions. The task becomes one of "catching ourselves in the act" of creating our own "reality" from the flow of events. Human truth is always an engagement of mind with experience.. We don't need to fear the collapse of our

personalistic belief system (the "box" we live in), nor our belief in absolute truth. **A strong desire to engage in the "quest for uncertainty" complements our anxiety that perhaps there is no absolute, objective ground to reality**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t-T_8)F_8(R_8(t-T_8) - N_8(t)M_8((R_8(t) \qquad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8} + \frac{N_8 dG_8}{dR_8} - \frac{dB_8}{dR_8} > \frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8} + \frac{dM_8}{dR_8}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_8 = \frac{N_8 G_8}{2M_8}\frac{d(8nF_8)}{dR_8} (\geq 0) \quad (8a)$$

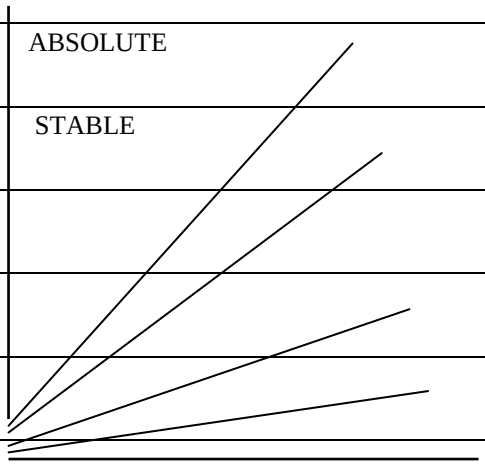
$$I_8 = -\frac{N_8 G_8}{2M_8}\frac{d(8nM_8}{dR_8} (\geq 0) \quad (8b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8 > L_8 - I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_8 > L_8 - I_8$, when $P_8^2/4 > I$ (8a)

$P_8 > L_8 - I_8 + \frac{(I_8 - P_8^2}{(L_8 - I)}$ when $P_8^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

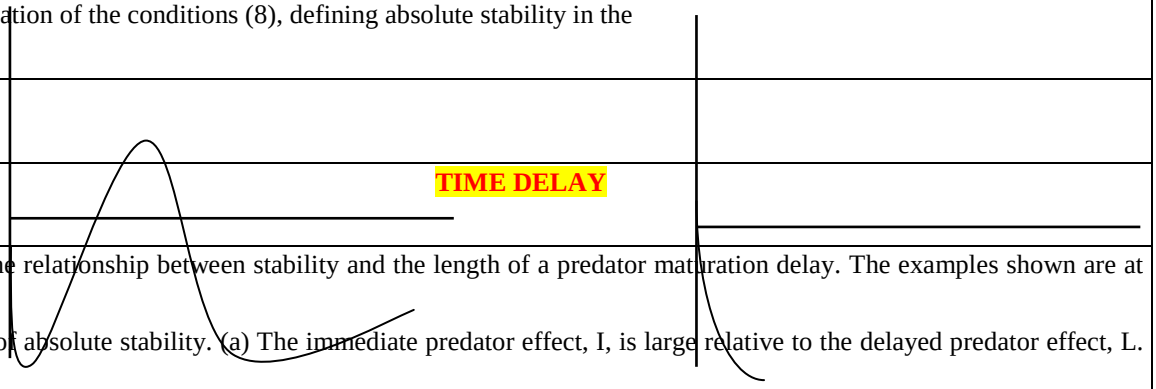



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this

rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior , we must note that the model could be applied to psychological systems too, albeit we are still not sure of the fact how the wave function collapses thereby due to perception.

Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility. He privatizes the sublime, having learned from the fate of his predecessors that the public can never be more than beautiful."(11) Notice is made here of what is indeed the telling feature. Levinas is praised, for example, for his alarm when "the social will [is] sought in an ideal of fusion . . . the subject losing himself in a collective representation, in a common ideal. . . . It is the collectivity which says 'us,' and which, turned toward the intelligible sun, toward the truth, experience, the other at his side and not face to face with him. The tellingly titled D'un ton apocalyptique contains passages that are particularly excited: These people situate themselves outside the ordinary, but they have in common this: they describe themselves as having an immediate and intuitive relation with mystery. And they want to attract, to seduce, and lead others to the mystery, through mystery. This agogic function of the leader of men, of the duce, of the Führer, of the leader, places him above the crowd that he manipulates with the aid of a small numbers of adepts joined together in a sect with a secret language, a clique or a small party with its

ritualized practices. The mystifiers pretend to have exclusive access to the privilege of a secret mystery. . . . The revelation or the unveiling of the secret is something that they jealously reserve for themselves. Jealousy is here a major characteristic. There is not narcissism and non-narcissism; there are narcissisms that are more or less comprehensive, generous, open, extended. What is called non-narcissism is in general but the economy of a much more welcoming, hospitable narcissism, one that is much more open to the experience of the other as other. I believe that without a movement of narcissistic reappropriation, the relation to the other would be absolutely destroyed; it would be destroyed in advance. The relation to the other--even if it remains asymmetrical, open, without reapporpriation--must trace a movement of reappropriation in the image of oneself for love to be possible, for example. Love is narcissistic. Beyond that, there are little narcissisms, there are big narcissisms. . . .(23) **Justice of the Pieces Deconstruction as a Social Psychology Douglas Collins** Sabine Wilke cogently argues for the rise of the aesthetic dimension in philosophical writing by showing how the metaphorical and poetic style of Heidegger, the principles of phrase configurations with Adorno, or the topographic writing of derrida replace through their very form something what traditionally would have been subject to discursive exposure The concern shifts from telling to showing, showing first and foremost that something remains untold. An exemplary book in this aesthetic drive is Avital Ronell's Telephone Book where the playful layout and explicit break with traditional design interrupt the sequential logic of the paragraphs thereby performing the very argument at stake: open up new circuits of meaning and signification. The engagement with form and publishing formats is only consequential since media theory as exposed by Walter J. Ong, Eric A. Havelock, Parry Miman or Marshall McLuhan besides many others made it clear that different communication technologies condition the mode of thought of a given culture The fractals of Benoit Mandelbrot or the butterfly shaped Lorenz attractor are the popular icons of this turn in the sciences. For communication studies and philosophy, Vilém Flusser was one of the pioneer thinkers reflecting the rise of what he called “techno-images.” The photograph is the paradigmatic case which condenses the whole crises of Western

thought and culture on its grainy surface. It subverts the linear logic of writing.6 Flusser's central diagnosis detects a shift from the grand narratives associated with the literate mind and Modernity towards creating and envisioning the world through image based practices. Writing himself in a succinct, visually inspired style, Flusser contents that the web of meaning, spun throughout the centuries by our textual practices, has become transparent. We discover the void between the letters through which the world escapes: "We are alienated from the world of text, precisely because we see through them as our own product."7 As a consequence we are confronted with a world of dissociated facts and information bits. Flusser hails this fragmented world as the advent of the zero dimensional point-universe that will provide the new raw material for our image based practices. Techno-images are nothing else than a conscious, inherently artistic drawing-together of point-elements as exemplified first and foremost by the silver-grains of the photograph but also by more recent pixel based techno-images. Similar to blueprints, designs or statistical curves, the weather map for example visualizes and contracts a tremendous amount of meteorological data from widely distributed stations into a single image of high- and low pressure

areas, temperatures, wind speeds, etc. Techno-images "compute" in this sense the universe of unimaginable, abstract information-bits into a concrete surface – the image. Echoing Paul Feyerabend, these image practices have to be understood as creative practices where the truth of scientific models has been replaced by its aesthetic quality **Aesthetics, Writing, Networked Computers Jörg Muller Dissertation Submitted to the Division of Media and Communications of The European Graduate School in Candidacy for the Degree of Doctor of Philosophy**

**Lothar Schäfer (1) and Sisir Roy (2) in** Quantum Reality, the Importance of Consciousness in the Universe, the Discovery of a Non-Empirical Realm of Physical Reality, and the Convergence with Ancient Traditions of Indian and Western Philosophy investigate metaphysical aspects that include 1) the discovery of a non-empirical part of physical reality in a realm of potentiality; 2) the emanation of the empirical world out of a realm of non-material forms; 3) the discovery that the nature of physical reality is that of an indivisible Wholeness – the One; and 4)

panpsychism: the possibility that the One is aware of its processes like a Cosmic Consciousness. The convergence of powerful traditions of seemingly disparate cultures is particularly important to point out in the present process of globalization, when a unifying view is needed to avoid controversy and conflict. To quote them in extenso : This is the nonlocality of the quantum world. Menas Kafatos and Robert Nadeau (1990) have drawn a remarkable conclusion from this phenomenon: If reality is nonlocal, the nature of the universe is that of an undivided wholeness. Because our consciousness has emerged from this wholeness and is part of it, it is possible to conclude that an element of consciousness is active in the universe: a Cosmic Consciousness. Pre-reflective self-consciousness is pre-reflective in the sense that (1) it is an awareness we have before we do any reflecting on our experience; (2) it is an implicit and first-order awareness rather than an explicit or higher-order form of self-consciousness. Indeed, an explicit reflective self-consciousness is possible only because there is a pre-reflective self-awareness that is an on-going and more primary self-consciousness. Although phenomenologists do not always agree on important questions about method, focus, or even whether there is an ego or self, they are in close to unanimous agreement about the idea that the experiential dimension always involves such an implicit pre-reflective self-awareness. In line with Edmund Husserl (1959, 189, 412), who maintains that consciousness always involves a self-appearance (Für-sich-selbst-erscheinens), and in agreement with Michel Henry (1963, 1965), who notes that experience is always self-manifesting, and with Maurice Merleau-Ponty who states that consciousness is always given to itself and that the word ‘consciousness’ has no meaning independently of this self-givenness (Merleau-Ponty 1945, 488), Jean-Paul Sartre writes that pre-reflective self-consciousness is not simply a quality added to the experience, an accessory; rather, it constitutes the very mode of being of the experience: This self-consciousness we ought to consider not as a new consciousness, but as the only mode of existence which is possible for a consciousness of something (Sartre 1943, 20 [1956, liv]). The notion of pre-reflective self-awareness is related to the idea that

experiences have a subjective 'feel' to them, a certain (phenomenal) quality of 'what it is like' or what it 'feels' like to have them.

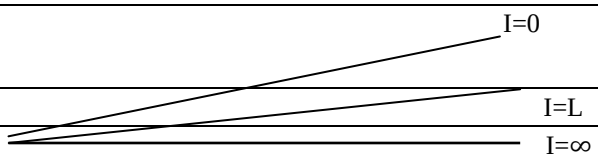
Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

TOTAL

ABSOLUTELY STABLE

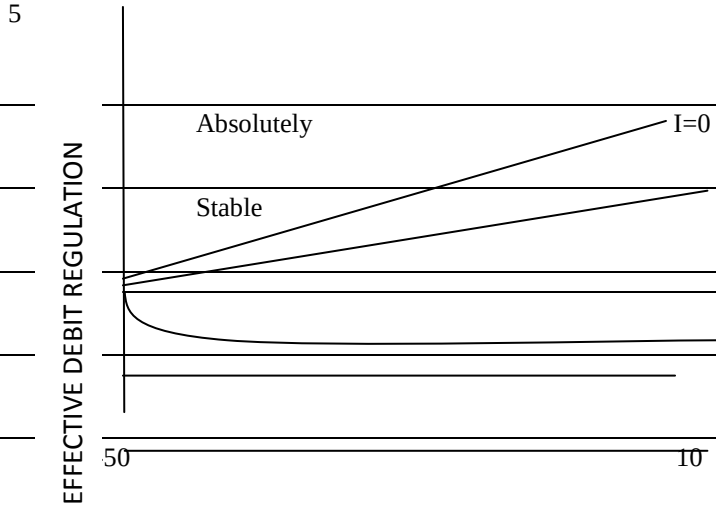



**DELAYED DEBIT EFFECT**

FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

## Utility Analysis And neural Network

Caroline Williams investigated affective density of the political and its effect on our understanding of political subjectivity. Taking up Spinoza's challenge to think about affect beyond corporeal embodiment, he argues that there is a modality of affectivity that cannot simply be inscribed within the borders of subjectivity. theorising affect as an impersonal force anchored in a relational ontology that gives due recognition to the circulation of affects, as well as to their ambivalent structure in creating sites of identification, and he utilises this ontology to reflect on the dynamic of the political and the shape of political subjectivity. Williams argues that Spinoza's philosophy (through ideas of conatus and imagination) offers the conceptual resources to reconfigure the composition of affective subjectivity as a transindividual social bond and as an unconscious dynamic of ethico-political existence. **Subjectivity (2010) 3, 245–262. doi:10.1057/sub.2010.15 Affective processes without a subject: Rethinking the relation between subjectivity and affect with Spinoza Caroline Williams** Phenomenological Time bears ample evidence and infallible observatory and apodictic evidence albeit with a sense of chagrined circumvention and defeated disgruntlment circularity of both psychological and phenomenological interpretations of art which are sometimes insightful, but most often work on the logic of cyclical time and unavoidable drives and representations played out as a ‘natural’ order: where ‘men can’t avoid being boys’ a circularity which valorizes a helplessness/truthfulness conundrum. One begins with a classical phenomenological fallacy with Duchamp’s Fountain, does the urinal keep returning or do we keep returning to the urinal? Put in this way, it is a male question, something that phenomenology conveniently conceals when positing the “I”, or indeed, the “we”. But we must be careful not to simply assume that the phenomenological first person is white, heterosexual, able-bodied and male, providing a foundation upon which knowledge should be based as a universal condition. Deleuze’s second critique of time is linked to Kant. This model of time extends not only to phenomenology but to the phenomenological interpretation

of art and is seen as a straight line. In the Critique of Pure Reason, Kant cuts up circular time and reassembles it as a series of sensory experiences which memory processes and which constitute subjectivity. Whereas the circular model of time (and interpretation) inscribes fatalism into the subject and the world, the phenomenological method is in danger of solipsism, where the subject constitutes the world. Much of embodied philosophy and enactivist theories of cognition go down this line with a plethora of art historical analyses by the likes of Rosalind Krauss, Michael Fried and others taking the lead from Merleau-Ponty. These approaches successfully challenge the **homunculus model of a Cartesian mind-body** dualism where there is a supervisor in the mind who processes experience and even consciousness and to whom all representations are directed. But phenomenology is mainly concerned with perceptual consciousness, the processing of colors, shapes, forms, motion in art and the world. The senses and sensorimotor processes are the fundamental pathways by which the interior is connected to the exterior in a chiasm of self-constitution in the world. The problem with this is that we explore the art through our sense perceptions, and what does art reveal to us?—our sense perceptions! Such solipsistic approaches have real problems relating complex and abstract conceptual production to bodily and sensorimotor contingencies, and they nearly always posit the subject as the site and object of discovery in the world. To be fair, phenomenology has moved on to rather more sophisticated models of intersubjectivity, taking into account the specificities of different kinds of bodies and ways of being beyond the white, male heterosexual as the common denominator, yet many naïve approaches to art perpetuate the notion that perceptual, rather than conceptual processes and evolutionary neural patterns are primary, fundamental ways to understand the art experience In fact, these aspects are only part of art production and reception, aspects which are often filtered out or made known to us in order for us to parody such trigger responses, as in much of Duchamp's work, which holds in contempt such retinal approaches to art. If we take the **subject of abstract expressionism** in the 1950s or later minimalist art, art history, teaching and interpretation still insists on putting the act of perceiving as fundamental to the constitution not only of the phenomenological first

person but also of art itself. The male act of usually standing up, which consists of directing a stream with the hand(s) into a container or on a flat surface, sand or snow in order to leave patterns and traces (the usual child's prank) underlies the logic of macho 1950s Abstract Expressionist culture, especially that of Jackson Pollock, so effectively parodied by Carolee Schneeman's Vagina Scroll, 1975. Here, painting or writing is a kind of a discharge which invests the **object and space with** the phenomenological "I", disciplined by existential struggle and reflected back by the traces, signatures and gestures left behind in the work of art, rediscoverable as an element in the making of other subjectivities.It is precisely this kind of solipsism which is blind to the art event as something well in excess of re-establishing habits of subjective self-identity through the production and interpretation of art. The attempt to reinforce a phenomenological identity upon the event by personalizing sense experience is unable to see art as something which remodels the phenomenological 'I' and neural plasticity itself. We learn new technologies and redraw our neural pathways through such new and spontaneous experiences which draw us into intersubjective worlds. Duchamp's Fountain and Deleuze's Repetition and Différence Grégory Minis sale (I have made some changes) The relation of different groups to deconstruction is settled by its relation to groups. An American enthusiast--one who has had a role in the development of our affection for Derrida--notices that the motivational chains are phobically driven by concerns with social scale. Richard Rorty: "He wants to figure out how to break with the temptation to identify himself with something big. . . ."(10) And: "So I take Derrida's importance to lie in his having had the courage to give up the attempt to unite the private and the public, to stop trying to bring together a quest for private autonomy and an attempt at public resonance and utility.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^\infty \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$then given the constraints of $f(T_8)$, $(A1.5) can\ never\ be\ satisfied\ Co$nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8 \lambda_8 e^{-T_8 \lambda_8} + [(U_8 - V_8) M_8 + Y_8] - [((U_8 - V_8) M_8 -] e^{-T_8 \lambda_8} = 0$$

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$\boldsymbol{Z}_8^2 + M_8 Z_8 \sin(Z_8 T_8) - (A_8 M_8 + Y_8) + (A_8 M_8 + X_8)\cos(Z_8 T_8) = 0$$

$$(A_8 M_8) Z_8 - M_8 Z_8 \cos(Z_8 T_8) + (A_8 M_8 - X_8)\sin(Z_8 T_8) = 0$$

Where $A_8 = U_8 - V_8$Using (A2.3), substitute for $\boldsymbol{Z}_8^2$ and $Z_8$in (A2.2):

$W_8^2 (X_8^2 - 3 M_8 A_8 X_8 - M_8^2 (X_8 - Y_8) + 2 M_8^3 A_8 + 2 A_8^2 M_8^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_8$=1-cos $(Z_8 T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8 / 2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_8 \le 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2 (X_8 - Y_8)$$

$$Q_8(2, A_8) = (A_8 + 2M_8)^2 (2A_8 M_8 - (X_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2\left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8\left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8\left(2X_8 + Y_8\right)\right)$$

$$+2A_8\left(X_8 + Y_8\right)$$

$$=2M_8^2 W_8^2 (M_8 + A_8) + M_8 A_8^2 W_8 + (2A_8 M_8 - X_8 + Y_8)(A_8 W_8 + M_8 W_8^2$$

$$+ (2 - W_8)((2M_8 X_8 + M_8 Y_8) W_8 + A_8 (X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $W_8$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in *A* will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range $0 \le W_8 \le 2$.Defining

$Q_8(W_8, A_8) = a_8 W_8^2 + c_8$ *where* $a_8$ $b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \square - b_8 \square 0$$

And

$$b_8^2 \square 4a_8 c_8$$

$$Case(i): Y_8 = 0$$

$Let\ Y_8 = 0$ And try the substitution A= $X_8 = /2M_8\, in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8 /$ $2M_8$) is positive for *W<2. Since* $Q_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\ X_8$ ,

$Let\ Y_8 = \ y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 + \ n_8 y_8)$ we have

$$W_8^2\, y_8 \left[X_8 n_8 \frac{(1 + n_8 y_8)}{(2 + n_8 y_8)} + M_8^2 \left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8 X_8 \left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2}+2y_8\left(n_8-1\right)+\frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)}\right]=0.\ (A2.12)$$

Since $n_8$ and $A_8$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8=(2+n_8 y_8)$, *the condition*$(A2.88?)$*can be written as*

$$\left[\frac{X_8(E_8-8)}{E_8^2}\right]^2\left[\frac{X_8^2}{M_8^2}-4y_8 E_8^2\frac{X_8}{M_8^2}+4y_8(n_8-2)E_8^3\right]$$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8/W_8^2)$ can be expressed as

$$2E_8\left[y_8E_8+(y_8E_8(8+y_8)(E_8-n_8))^{8/2}\right]>\frac{X_8}{M_8^2}>2E_8\left[y_8E_8-y_8E_8(8+y_8)(E_8-n_8))^{8/2}\right]$$

Thoughts on the Social and Political Implications of Correlationism .Lion Mirror Truth be told, as my thought has evolved the issue of correlationism had fallen off the radar for me. Somehow the debate had come to seem too "philosophical" to me, too "scholastic", too remote from what interests me: understanding why social assemblages are organized as they are, how power functions in social assemblages, and what we might do to address that power and change things. Somehow the question of whether or not we can get out of the correlation between thinking and being just came to seem remote from these sorts of issues. Somehow it seemed too epistemological. ASIDE: Numerous discussions over the years have led me to believe that the debate over correlationism is poorly understood (or maybe I just don't understand it). On countless occasions I've heard people say "of course we must relate to things in order to know them." Well yeah, of course! I don't think this is what the critic of correlationism is getting at. It seems to me that correlationism is something more robust than the theses that we must relate to something to know it.

Correlationism instead seems to require the theses that thought and being are indiscernible. Put more concretely, the correlationist is someone who argues that we either a) can never tell whether being is merely a construction of our thought (weak correlationism), or b) who argues that thought actually constructs being (strong correlationism). In other words, correlationism is another name for idealism. One can hold that we must relate to something in order to know it without being a correlationist. As an aside I should also add that I am a correlationist about some things. For example, I think money is something constructed by society and is therefore a strong correlationist when it comes to money. At any rate, for a long time I'd become rather indifferent to debates about correlationism and philosophies of access. I had learned the lessons of speculative realism– which I could have also learned, I think, from Deleuze and Guattari, the new materialist feminists, actor-network theorists such as Latour –and had moved on. However, occasionally you come across a tone of phrase that pitches something in a different light. In The Cut of the Real, Katerina Kolozova writes, …the political problem of contemporary philosophy identified by the 'new realists' is, in fact, the product of a more fundamental epistemic problem. In his book After Finitude, Quentin Meillassoux calls this problem 'correlationism' and identifies it as an essentially post-Kantian legacy, which continues to dominate and limit philosophy. As a matter of fact, correlationism lies at (eb) the heart of postmodern theory and consists in the premise that thought can only 'think itself,' that the real is inaccessible to knowledge and human subjectivity, and that there is nothing but discursive constructs that fully determine thinking and that are meth0dologically accounted for all the way down. (1 – 2) Thought thinking only itself: Thought only encountering itself. In the jargon of postmodern and poststructuralist lingo, this would be the thesis of infinite semiosis, where signs ("thoughts") only ever relate to other signs. Within this framework, discursivity comes to be the hegemonic framework defining all of being. At the level of politics and social theory more generally, if the correlationist thesis is true the

consequences are clear: all social phenomena are discursive and all solutions to social and political problems will be discursive. The sole sphere of the political will be the discursive and all questions of politics will be questions of speech-acts and interpretation. The problem here is not that many theorists recognize that the discursive and semiotic plays an important role in the social and the political. It does and I've repeated this tirelessly. The problem is with what happens when thought or the semiotic becomes a hegemon, an "all", foreclosing our ability to recognize other forms of power. What I've wanted to say is that not all power functions discursively. In my last post and elsewhere I spoke of some other forms of politics: Thermopolitics: The politics surrounding **energy in the form of calories and fuels such as gasoline** and coal, and how our life and our very bodies are structured by energy dependencies and by being trapped in particular distributive networks that render these forms of energy available. I'm being quite literal when I speak of energy, talking about the effects, for example, of the absence of food in certain educational environments on cognition, for example; and am generally hostile to metaphorical extensions of the concept of energy which I see as erasing the dimension of real materiality. Geopolitics: The role that features of natural and built geography such as mountain ranges, rivers, oceans, soil conditions, roads, housing design, etc., play in the form that social relations take and how they impact individual bodies. Chronopolitics: The way in which the structurations of time organize what is possible for us. For example, the structuration of the working day, how much we can say and comprehend at any given time, the impact of things like the invention of the clock, etc. **Oikopolitics: This would be the domain of political economy described so well by Marxists**. So five different types of politics: Semiopolitics (or what currently dominates critical theory), thermopolitics, geopolitics, chronopolitics, and oikopolitics. No doubt there are other sites of the political or political struggle that we could speak of, but this is a good start. Also, it should be obvious that these aren't exclusive domains, but are entangled in all sorts of important

ways. For example, something might take place at the level of semiopolitics (speech, law, rhetoric, norms, and communication) that has all sorts of impact at the level of thermopolitics. Congress might decide to cut programs that fund school meal programs. This, in turn, will have a thermodynamic impact on those students that go without the calories they need developmentally and cognitively to function in a particular way. There is an entanglement here of semiopolitical and thermopolitical domains. The young student here has been constrained both at the level of **semiotic phenomena and thermodynamic structures**. The point is that if true, semiotic intervention (speech-acts, protests, interpretations, deconstructions, etc) will not be an appropriate response to all political problems because social formations are not entirely structured by the semiotic. The child in that school does not suffer from a lack of the right signs, but from a lack of calories needed to run the engine of his thought and body. Certainly semiotic interventions might be needed to render that energy available, but it is the energy itself that is at issue and the absence of that energy that forms the spider web entangling him in his position. A correlationist perspective tends to erase this as even being a site of the political. **Larval Subjects January 28, 2014 Ontological Anarché Posted by larval subjects under Uncategorized January 27, 2014 Roden on Pluralism Over at enemy industry Roden has an excellent post up on the pluralism discussion. Check it out here. January 25, 2014 Different Senses of Pluralism and Ontology Posted by larval subjects under Uncategorized Responding to one of my comments over at his blog, Jeremy Trombley:** I have deleted restricted, reformulated some sentences .Kindly bear with me. Intention is to make the discursive enucleation consistent with the subject in question**.** Recourse to epistemology of the quantum mechanics generic and fervor of string theory is a sine qua non for the further comprehension of the theory and ontology .How would the world appear to us if its ontology was that of classical mechanics but every agent faced a restriction on how much they could come to know about the classical state? We show that in most respects

it would appear to us as quantum. The statistical theory of classical mechanics, which specifies how probability distributions over **phase space evolve** under Hamiltonian evolution and under measurements, is typically called Liouville mechanics, so the theory we explore here is Liouville mechanics with an epistemic restriction. The particular epistemic restriction we posit as our foundational postulate specifies two constraints. The first constraint is a classical analog of Heisenberg's uncertainty principle; the second-order moments of position and momentum defined by the phase-space distribution that characterizes an agent's knowledge are required to satisfy the same constraints as are satisfied by the moments of position and momentum observables for a quantum state. The second constraint is that the distribution should have maximal entropy for the given moments. Starting from this postulate, authors derive the allowed preparations, measurements, and transformations and demonstrate that they are isomorphic to those allowed in Gaussian quantum mechanics and generate the same experimental statistics. They argue that this reconstruction of Gaussian quantum mechanics constitutes additional evidence in favor of a research program wherein quantum states are interpreted as states of incomplete knowledge and that the phenomena that do not arise in Gaussian quantum mechanics provide the best clues for how one might reconstruct the full quantum theory. I: http://dx.doi.org/10.1103/PhysRevA.86.012103 **Reconstruction of Gaussian quantum mechanics from Liouville mechanics with an epistemic restriction Phys. Rev. A 86, 012103 – Published 10 July 2012 Stephen D. Bartlett, Terry Rudolph, and Robert W. Spekkens.** Deleuze in fact characterizes modes of existence, with their powers and capacities. The answer is this: Deleuze approaches modes of existence, ethically speaking, not in terms of their will, or their conscious decision making power (as in Kant), nor in terms of their interests (as in Marx, for example), but rather in terms of their drives. For Deleuze, conscious will and preconscious interest are both subsequent to our unconscious drives, and it is at the level of the drives that we have to aim our ethical analysis. Here, I would like to

focus on two sets of texts on the drives taken, not from Nietzsche and Spinoza, but rather from Nietzsche and Leibniz (Leibniz being one of the first philosophers in the history of philosophy to have developed a theory of the unconscious). The first set of texts comes from Nietzsche's great early book entitled Daybreak, published in July 1881. Nietzsche first approaches the question of the drives by giving us an everyday scenario: "Suppose we were in the market place one day," he writes, "and we noticed someone laughing at us as we went by: this event will signify this or that to us according to whether this or that drive happens at that moment to be at its height in us—and it will be a quite different event according to the kind of person we are. One person will absorb it like a drop of rain, another will shake it from him like an insect, another will try to pick a quarrel, another will examine his clothing to see if there is anything about it that might give rise to laughter, another will be led to reflect on the nature of laughter as such, another will be glad to have involuntarily augmented the amount of cheerfulness and sunshine in the world—and in each case, a drive has gratified itself, whether it be the drive to annoyance, or to combativeness or to reflection or to benevolence. This drive seized the event as its prey. Why precisely this one? Because, thirsty and hungry, it was lying in wait" (D 119). This is the source of Nietzsche's doctrine of perspectivism ("there are no facts, only interpretations"), but what is often overlooked is that, for Nietzsche, it is our drives that interpret the world, that are perspectival—and not our egos, not our conscious opinions. It is not so much that I have a different perspective on the world than you; it is rather that each of us has multiple perspectives on the world because of the multiplicity of our drives—drives that are often contradictory among themselves. "Within ourselves," Nietzsche writes, "we can be egoistic or altruistic, hard-hearted, magnanimous, just, lenient, insincere, can cause pain or give pleasure" (Parkes, pp. 291-292). We all contain such "a vast confusion of contradictory drives" (WP 259) that we are, as Nietzsche liked to say, multiplicities, and not unities. Moreover, these drives are in a constant struggle or combat with each other:

my drive to smoke and get my nicotine rush is in combat with (but also coexistent with) my drive to quit. This is where Nietzsche first developed his concept of the will to power—at the level of the drives. "Every drive is a kind of lust to rule," he writes, "each one has its perspective that it would like to compel all the other drives to accept as a norm" (WP 481). DELEUZE AND THE QUESTION OF DESIRE: TOWARD AN IMMANENT THEORY OF ETHICS Daniel W. Smith. **In on a Possible Physical Metatheory of Consciousness Miroljub Dugic et al** show that the modern quantum mechanics, and particularly the theory of decoherence, allows formulating a sort of a physical metatheory of consciousness. Particularly, the analysis of the necessary conditions for the occurrence of decoherence, along with the hypothesis that consciousness bears (more-or-less) well definable physical origin, leads to a wider physical picture naturally involving consciousness. This can be considered as a sort of a psycho-physical parallelism, but on very wide scales bearing some cosmological relevance. Open Systems & Information Dynamics Volume 9, Number 2, 153 (2002) arXiv: quant-ph/0212128. On a broader perspective and background, the theoretical computation of the universe is justified.IN physics and cosmology, digital physics is a collection of theoretical perspectives based on the premise that the universe is, at heart, describable by information, and is therefore computable. Therefore, according to this theory, the universe can be conceived of as either the output of a deterministic or probabilistic computer program, a vast, digital computation device, or mathematically isomorphic to such a device. Digital physics is grounded in one or more of the following hypotheses; listed in order of decreasing strength. The universe or reality: is essentially informational (although not every informational ontology needs to be digital); is essentially computable (the pancomputationalist position); can be described digitally; is in essence digital; is itself a computer (pancomputationalism); is the output of a simulated reality exercise. Following Jaynes and Weizsäcker, the physicist John Archibald Wheeler wrote the following: [...] it is not unreasonable to imagine that

information sits at the core of physics, just as it sits at the core of a computer. (John Archibald Wheeler 1998: 340) It from bit. Otherwise put, every 'it'—every particle, every field of force, even the space-time continuum itself—derives its function, its meaning, its very existence entirely—even if in some contexts indirectly—from the apparatus-elicited answers to yes-or-no questions, binary choices, bits. 'It from bit' symbolizes the idea that every item of the physical world has at bottom—a very deep bottom, in most instances—an immaterial source and explanation; that which we call reality arises in the last analysis from the posing of yes–no questions and the registering of equipment-evoked responses; in short, that all things physical are information-theoretic in origin and that this is a participatory universe. (John Archibald Wheeler 1990: 5) David Chalmers of the Australian National University summarised Wheeler's views as follows: Wheeler (1990) has suggested that information is fundamental to the physics of the universe. According to this 'it from bit' doctrine, the laws of physics can be cast in terms of information, postulating different states that give rise to different effects without actually saying what those states are. It is only their position in an information space that counts. If so, then information is a natural candidate to also play a role in a fundamental theory of consciousness. We are led to a conception of the world on which information is truly fundamental, and on which it has two basic aspects, corresponding to the physical and the phenomenal features of the world. (Wikipedia). **Participatory consciousness forms the bastion, pillar, post, stylobate and fulcrum of the redressal measures and corrective steps taken for circumvention of many problems, identified and unidentified, spread widest commonalty. Anne Baring's sententious and pithy article will be good to have a timely outlook.** I would like to open my contribution to this meeting with the words of Ken Wilber that I quoted in the first Millennium Symposium: 'The secret impulse of life is towards greater consciousness. Maybe the evolutionary sequence is from matter to body, to mind, to soul, to spirit - each transcending and including, each with a greater depth and greater

consciousness and wider embrace' Ken Wilber, and with these from Rupert Sheldrake's contribution to the same meeting: "It is possible to make quite a good case that the sun could be conscious and if the sun, then why not the stars? Why not the galaxies? Why should there not be galactic minds? And if they communicate with each other then the distances are far too great for light to be the medium of transmission. We are into inter-galactic telepathy." After many decades of study, I know of only one diagram in all religious traditions that can adequately explain these statements; that can give us both a blue-print of the many levels of consciousness and a route-map to help us reach beyond our present level of understanding. I would like to introduce a diagram of the four worlds of Kabbalah to illustrate the theme that spirit has brought these four worlds into being and that we, living in the fourth and manifest world know nothing of the existence of the other three invisible ones. It is the clearest diagram I know of that can explain to us the multi-levelled structure of consciousness and our place in the great chain of being. We are indeed, as Wilbur suggests, evolving from matter to body, to mind, to soul, to spirit but where have we come from? It is possible that, as participants in the emanation of spirit into manifest life, we have also come from spirit to soul, to mind, to body, to matter. Everything, from the immense galactic energies to the minute forms of matter participates in one life, is generated from one creative spirit. There is nothing outside this life, this spirit. This diagram suggests that truth is not an article of faith but the experience of participating in the undiscovered reality of what has brought us into being. We cannot know truth until we have entered into communion with that reality. Continuing in similar vein **Michel Bitbol dilates upon the dialectic deliberation and polemical argumentation between** hermeneutists and eliminativists. When he formulated the program of Neurophenomenology, Francisco Varela suggested a balanced methodological dissolution of the “hard problem” of consciousness. He shows that his dissolution is a paradigm which imposes itself onto seemingly opposite views, including materialist approaches. I also point out that

Varela's revolutionary epistemological ideas are gaining wider acceptance as a side effect of a recent controversy between hermeneutists and eliminativists. Finally, he emphasizes a structural parallel between the science of consciousness and the distinctive features of quantum mechanics. This parallel, together with the former convergences, point towards the common origin of the main puzzles of both quantum mechanics and the philosophy of mind: neglect of the constitutive blindspot of objective knowledge. Phenomenology and the Cognitive Sciences Phenomenology and the Cognitive Sciences **Phenomenology and the Cognitive Sciences 2002, Volume 1, Issue 2, pp 181-224 Science as if situation mattered Michel Bitbol**

## ABSOLUTE STABILITY

Hence if $X_8 / M_8^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_8$ ) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8 - 8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2 - 2y_8\frac{(E_8 + n_8^2 y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right) - 4y_8(E_8 - n_8)$$

And has the solution

$$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8 - 8)}[y_8(E_8 + n_8^2 y_8) + [y_8^2 (E_8 + n_8^2 y_8)^2 + 4 y_8 (E_8 - 8)(E_8 - n_8)]^{8/2}]$$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2E_8[y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

**ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION**

**CONSOILIDATION**

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + (M_8 + U_8)\lambda_8 - (M_8 + V_8)\lambda_8 e^{2T_8\lambda_8} + (M_8 U_8 + Y_8)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8\, e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + (M_8 + V_8) Z_8 \sin(Z_8 T_8) - (U_8 M_8 + Y_8) + (M_8(U_8 + V_8) - X_8)$$

$$\mathrm{X}\cos(Z_8 T_8) - 2M_8 V_8 \cos^2(Z_8 T_8) + M_8 V_8 = 0$$

$$(M_8 + U_8) Z_8 - (M_8 + V_8) Z_8 \cos(Z_8 T_8) + (M_8(U_8 + V_8) - X)$$

$$\mathrm{X}\sin((Z_8 T_8) - M_8 V_8 \sin(Z_8 T_8) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 =$

$1 - \cos Z_8\, T_8$and $A_8 = U_8 - V_8$we obtain

$$4M_8^2 V_8^2 W_8^4 + 2M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8(M_8 + V_8) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2[M_8 - V_8)^2(2A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8(M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8(A_8^2 + 4(X_8 + Y_8))]$$

$$+W_8\left[A_8^2+2X_8\right)\left(2A_8M_8-X_8\right)+2A_8Y_8(M_8+V_8)-M_8A_8^3]+A_8^2\left(X_8M_8-Y_8\right)=0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8>0$$

$$and$$

$$C_8>(X_8-Y_8)/2M_8$$

$recalling\ that\ A_8$= $U_8$and defining $C_8=U_8+V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8(0,U_8V_8=A_8^2(X_8+Y_8)$$

$$Q_8\left(2,U_8V_8\right)=\left(C_8+2M_8^2\right)^2(2C_8M_8+Y_8-X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0$W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$the existence of a D-position at this point is determined by letting cos $(Z_8T_8)=-1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all $T$, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**ABSOLUTE STABILITY IN PREDATOR- PREY MODELS**

**APPLIED TO ACCETUATION ATTRITION MODELS**

Equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability

criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The Percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending

upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1985); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

**DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY**

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 8985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_8}{dt} = [N(t)]_8[(F_8)(R(t)]_8 - [M]_8[R(t)]_8$$

$$\frac{d[R(t)]_8}{dt} = [B]_8\,[(R)_8(t-T_8)] - [D]_8[(R(t)]_8 - [N(t)]_8\left[G\big(R(t)\big)_8\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_8}{dt} = a_8[N(t)]_8 - a_8[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_8 = a_3[N(t)]_3 - a_4[N(t)]_8$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_8 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.8b, which represents a system with self-regulation acting through recruitment $(dD/dR<0)$. The

regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 8a is at the

**TIME DELAY**

FIG. 8. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a dominant eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.8$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

DELAY DUE TO CONGREGATIONAL AND CONFEDERATIONAL DEBIT CREDIT TYEPES

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_8}{dt} = N_8(t-T_8)F_8(R_8(t-T_8) - N_8(t)M_8((R_8(t) \quad (5)$$

$$\frac{dR_8(t)}{dt} = B_8\big(R_8(t) - D_8(R_8(t)\big) - N_8(t)G_8(R_8(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_8}{dR_8} + \frac{N_8 dG_8}{dR_8} - \frac{dB_8}{dR_8} > \frac{N_8 G_8}{2M_8}\left[\frac{dF_8}{dR_8} + \frac{dM_8}{dR_8}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_8/dR_8$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_8$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_8$) and the other derived from the non-delayed or "immediate" processes($I_8$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

## ABSOLUTE STABILITY

$$L_8 = \frac{N_8 G_8}{2M_8}\frac{d(8nF_8)}{dR_8} (\geq 0) \quad (8a)$$

$$I_8 = -\frac{N_8 G_8}{2M_8}\frac{d(8nM_8}{dR_8} (\geq 0) \quad (8b)$$

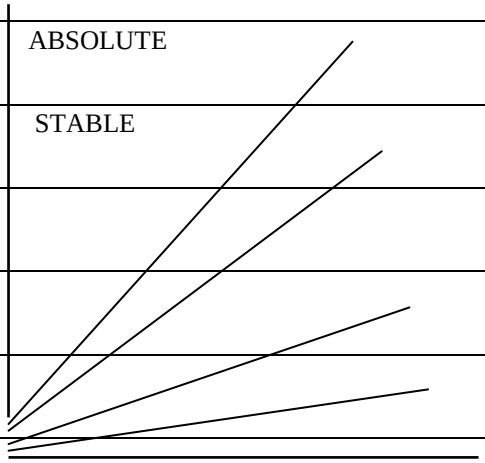
 

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_8 > L_8 - I_8$, must be modified ahe $dM_8/dR_8$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_8 > L_8 - I_8,$ when $P_8^2/4 > I$ (8a)

$P_8 > L_8 - I_8 + \frac{(I_8 - P_8^2}{(L_8 - I)}$ when $P_8^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit

(like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

See the diagram Nunney gives for predator and prey analysis. We can simulate the same for the financial system in the foregoing too.

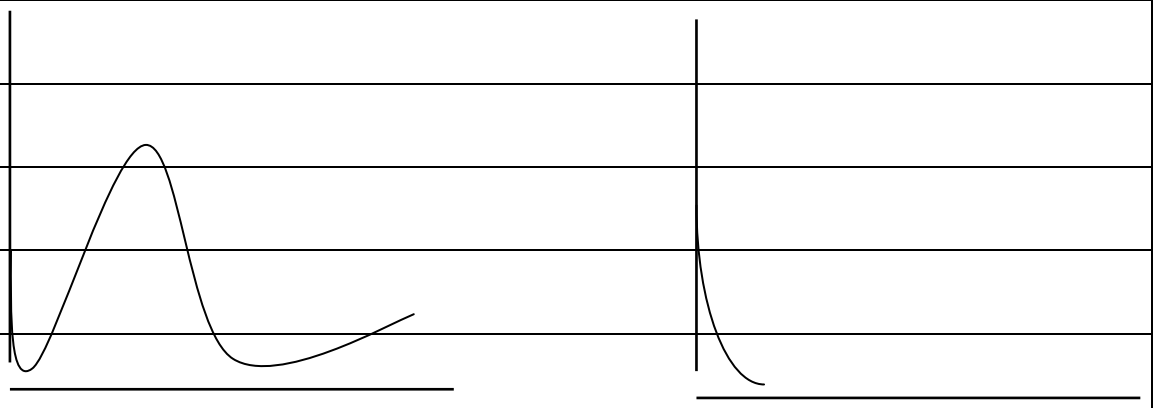
Can never completely offset the lagged effect,$L_8$; condition (8b) ensures that the total debit regulation, $P_8$, must always be positive even when $I_8$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_8$ and $L_8$. As $I_8$ increases relative to $L_8$ the zone of stability is enlarged; however,even when $I_8$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_8$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (8). Parameters in each case M=8; $dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8 = 0.05$ For (a) $NGdF/dR = 0.3, -N_8G_8dM_8/dR_8 = 0.6$; for (b)$NG_8dM_8/DR_8 = 0$

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_8t}{dt} = N_8(t - T_8)F_8\big(R_8(t - T_8)\big) - N_8(t)M_8\big(R_8(t)\big)$$

(9)

$$\frac{dR_8t}{dt} = B_8\big(R_8(t - T_8)\big) - D_8\big(R_8(t)\big) - N_8(t)G_8\big(R_8(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) and the effective prey regulation ($dD_8/dR_8 + NdG_8/dR_8 - dB_8/dR_8$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_8/dR_8$ is a measure of the control acting through the debit delay and setting $dB_8/dR_8 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_8/dR_8$ is a measure of the control acting through the credit delay although setting $dF_8/dR_8 = 0$ does not negate the effect of a predator delay; the value of $M_8$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the

debit delay, as measured by the magnitude of $dF_8/dR_8$and normalized by $M_8$ (see8a). The absolute stability of such a system is summarized in Fig. 8 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

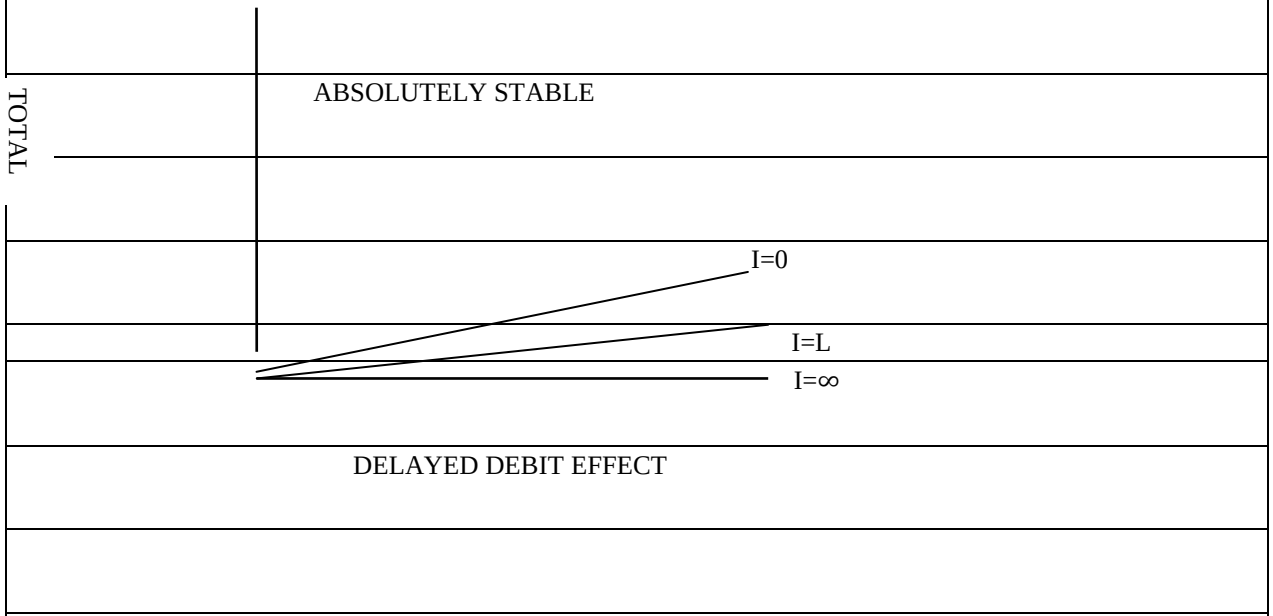
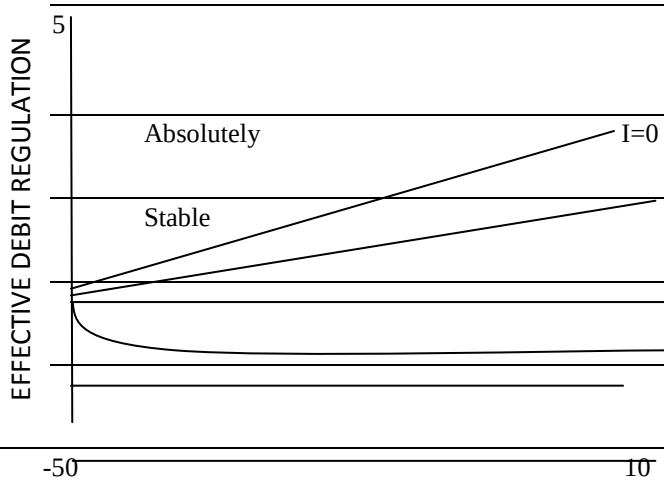



FIG8. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_8/dR_8 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison?

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.




DELAYED CREDIT EFFECT (lack of balance)

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_8/dR_8 + N_8 dG_8/dR_8 + dB_8/dR_8)/M_8$). The immediate ($I_8$) and lagged ($L_8$) predator effect are

defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_8$(see 8b), is increased relative to the lagged part, $L_8$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_8 > I_8$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_8$is somewhat large than $I_8$ then the condition constraining prey regulation of recruitment ($dB_8/dR_8 < 0$) is

$$\frac{dD_8}{dR_8} + N_8\frac{dG_8}{dR_8} + \frac{dB_8}{dR_8} > \frac{(L_8 - I_8)M_8}{2} \qquad (10)$$

If $L_8 < I_8$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_8$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_8$

The - $D_8$ positions are defined by letting $\lambda_8 = iZ_8$ and separating (A1.2) into real and imaginary parts:

$$Z_8^2 + V_8 Z_8 \int_0^\infty \sin(Z_8 T_8) f(T_8) dT_8 - (X_8 + Y_8) = 0$$

$$Z_8 \left( U_8 - V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8 \right) = 0$$

Noting that $X_8 + Y_8 = 0$ is we assume X+Y=0.It follows that $Z_8 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_8 = V_8 \int_0^\infty cos\left(Z_8 T_8\right) f\left(T_8\right) dT_8$$

If $U_8 > |V_8|$ then given the constraints of $f(T_8)$, $(A1.5) can\ never\ be\ satisfied$ Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A8.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_8^2 + (U_8 - V_8 + M_8)\lambda_8 - M_8 \lambda_8 e^{-T_8 \lambda_8} + [(U_8 - V_8) M_8 + Y_8] - [((U_8 -$$

$$V_8 M_8 -] e^{-T_8 \lambda_8} = 0$$

(For definitions see Appendix8) We follow the technique outlined in Appendix 8 and let $\lambda_8 = iZ_8$ and separate (A2.8) into real and imaginary parts

$$\boldsymbol{Z}_8^2 + M_8 Z_8 \sin(Z_8 T_8) - (A_8 M_8 + Y_8) + (A_8 M_8 + X_8)\cos(Z_8 T_8) = 0$$

$$(A_8 M_8) Z_8 - M_8 Z_8 \cos(Z_8 T_8) + (A_8 M_8 - X_8)\sin(Z_8 T_8) = 0$$

Where $A_8 = U_8 - V_8$Using (A2.3), substitute for $\boldsymbol{Z}_8^2$ and $Z_8$in (A2.2):

$W_8^2 (X_8^2 - 3 M_8 A_8 X_8 - M_8^2 (X_8 - Y_8) + 2 M_8^3 A_8 + 2 A_8^2 M_8^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_8$=1-cos $(Z_8 T_8)$.

In the absence of a delay the system is stable for $A_8 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_8 > (X_8 - Y_8)/2\ M_8$ Thus a necessary condition for absolute stability is

$A_8 > [\ X_8 - Y_8 / 2\ M_8$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_8$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_8$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_8, A_8) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_8 \le 2$ *is* of interest First consider the extreme values

$$Q_8(0, A_8) = A_8^2\ (\mathrm{X}_8 - Y_8)$$

$$Q_8(2, A_8) = \left(A_8 + 2\,\mathrm{M}_8\right)^2 (2A_8\,\mathrm{M}_8 - (\mathrm{X}_8 - Y_8))$$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_8}{\partial A_8} = W_8^2\left(2M_8^3 + 4A_8 M_8^2 - 3M_8 X_8\right) + W_8\left(3M_8 A_8^2 - 2A_8 X_8 + 2M_8\left(2X_8 + Y_8\right)\right)$$

$$+2A_8\left(X_8 + Y_8\right)$$

$$=2\,M_8^2 W_8^2 (M_8 + \mathrm{A}_8) + M_8\, A_8^2 W_8 + \left(2A_8 M_8 - X_8 + Y_8\right)(A_8 W_8 + M_8 W_8^2$$

$$+ (2- W_8)((2\,M_8\,X_8 + M_8\,Y_8)\,W_8 + A_8\,(X_8 + Y_8)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_8$quadratic in $\mathrm{W}_8$touches the $W$ –axis only at its minimum value in which case (A2.9) shows that further increases in *A* will life the function away from the axis and maintain absolute stability or when the function $Q_8$has no minimum in the range $0 \le W_8 \le 2$.Defining

$Q_8\left(W_8, A_8\right) = a_8 W_8^2 + c_8$ *where* $a_8\, b_8$ *and* $c_8$ Are functions of $A_8$we can define the condition for a minimum in the range 0-2

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_8 \square b_8 \square 0$$

And

$$b_8^2 \square 4a_8c_8$$

$$Case(i): Y_8 = 0$$

$Let\ Y_8 = 0$ And try the substitution A= $X_8 = /2M_8 in$ (A2.4) this leads directly to a linear equation with the solution

$W_8 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_8 X_8 /$ $2M_8$) is positive for *W<2. Since* $Q_8$ increases with $A_8$(see A2.9), increasing $A_8$will generate no further solutions of $Q_8$=0. Hence the condition A>X/2M defines absolute stability when $Y_8 = 0$.

$Case\ (ii): Y_8 = y_8\ X_8$ ,

$Let\ Y_8 =\ y_8\ X_8$And try the substitution $A_8 = X_8/(2 + A_8\ y_8)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_8/(2 +\ n_8 y_8)$ we have

$$W_8^2\ y_8\left[X_8 n_8 \frac{(1 + n_8 y_8)}{(2 + n_8 y_8)} + M_8^2\left(2 - n_8 + n_8 y_8\right)\right]$$

$$-W_8 X_8 \left[\frac{X_8(1+n_8 y_8)}{M_8^2(2+n_8 y_8)^2}+2y_8\left(n_8-1\right)+\frac{X_8^2(8+y_8)}{M_8^2(2+n_8 y_8)}=0.(A2.12)\right.$$

Since $n_8$ and $A_8$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_8$

Defining $E_8=(2+n_8 y_8), the\ condition (A2.88?) can\ be\ written\ as$

$[\frac{X_8(E_8-8)}{E_8^2}]^2[\frac{X_8^2}{M_8^2}-4y_8 E_8^2\frac{X_8}{M_8^2}+4y_8(n_8-2)E_8^3]$ 0.

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_8/W_8^2)$ can be expressed as

$2E_8[y_8E_8+(y_8 E_8(8+y_8)(E_8-n_8))^{8/2}]>\frac{X_8}{M_8^2}>2E_8[y_8E_8-y_8E_8(8+y_8)(E_8-n_8))^{8/2}]$

**ABSOLUTE STABILITY**

Hence if $X_8/M_8^2)$ is large or small then (A2.6) will have a real solution although if may not be in the range (0 $W_8$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_8-8)}{E_8^2}\left(\frac{X_8}{M_8^2}\right)^2-2y_8\frac{(E_8+n_8^2y_8)}{E_8}\left(\frac{X_8}{M_8^2}\right)-4y_8(E_8-n_8)$$ 0

And has the solution

$\frac{X_8}{M_8^2} > \frac{E_8}{(E_8-8)}[y_8(E_8 + n_8^2 y_8) + [y_8^2 (E_8 + n_8^2 y_8)^2 + 4 y_8 (E_8 - 8)(E_8 - n_8)]^{8/2}]$

Given $E_8 > n_8$ this lower limit to $X_8 / M_8^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_8 / M_8^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_8$ in the appropriate range and this limit is further lowered by (A2.84)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_8 < n_8$ $(ie., 2A_8 M_8^2 < X_8 - Y_8)$ then (A2.8)and (A2.8)show that $Q_8$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_8}{M_8^2} > 2 E_8 [y_8 E_8 - (y_8 E_8 (8 + y_8)(E_8 - n_8)]^{8/2}]$$

And

$$E_8 > n_8$$

Replacing $E_8$ and $n_8$ in these conditions leads to the form

If $A_8^2 > 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.

If $A_8^2 < 2Y_8$ then $A_8 > (X_8 - Y_8)/2M_8$.+ $(Y_8 - A_8^2/2)^2 /2M_8 (X_8 - Y_8)$.

Note that by normalizing $A_8$ by $M_8$ and $X_8$ and $Y_8$ by $W_8^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_8^2 + (M_8 + U_8)\lambda_8 - (M_8 + V_8)\lambda_8 e^{2T_8\lambda_8} + (M_8 U_8 + Y_8)$$

$$-[M_8(U_8 + V_8) - X_8]\, e^{2T_8\lambda_8} + M_8 V_8\, e^{2T_8\lambda_8} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_8 = \mathrm{i}\, Z_8$ *and* separate

(A3.1)into real and imaginary parts

$$Z_8^2 + (M_8 + V_8) Z_8 \sin(Z_8 T_8) - (U_8 M_8 + Y_8) + (M_8(U_8 + V_8) - X_8)$$

$$\mathrm{X}\cos(Z_8 T_8) - 2M_8 V_8 \cos^2(Z_8 T_8) + M_8 V_8 = 0$$

$$(M_8 + U_8) Z_8 - (M_8 + V_8) Z_8 \cos(Z_8 T_8) + (M_8(U_8 + V_8) - X)$$

$$\mathrm{X}\sin((Z_8 T_8) - M_8 V_8 \sin(Z_8 T_8) = 0$$

Using (A.3.3) to eliminate $Z_8$and $Z_8^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_8 = 1 - \cos Z_8 T_8$and $A_8 = U_8 - V_8$we obtain

$$4M_8^2 V_8^2 W_8^4 + 2M_8 V_8 W_8^3[(M_8 - V_8)^2 + A_8(M_8 + V_8) + 2(A_8 M_8 - X_8)]$$

$$+ W_8^2[M_8 - V_8)^2(2A_8 M_8 + Y_8 - X_8) + (2A_8 M_8 - X_8)^2 - A_8(M_8 - V_8)(2A_8 M_8 - X_8)$$

$$+ M_8 V_8(A_8^2 + 4(X_8 + Y_8))]$$

$$+ W_8[A_8^2 + 2X_8)(2A_8 M_8 - X_8) + 2A_8 Y_8(M_8 + V_8) - M_8 A_8^3] + A_8^2(X_8 M_8 - Y_8) = 0$$

Being a quartic in $W_8$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_8 > 0$$

$$and$$

$$C_8 > (X_8 - Y_8)/2M_8$$

$recalling\ that\ A_8 = U_8$ and defining $C_8 = U_8 + V_8$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_8(0, U_8 V_8 = A_8^2 (X_8 + Y_8)$$

$$Q_8\left(2, U_8 V_8\right) = \left(C_8 + 2M_8^2\right)^2 (2C_8 M_8 + Y_8 - X_8)$$

So that given (A3.6) the extreme $W_8$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_8$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_8$the existence of a D-position at this point is determined by letting cos $(Z_8 T_8) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 e^{-\tau_8\lambda_8} + X_8 + Y_8 = 0 \qquad (A1.1)$$

Where $V_8 = dB_8/dR_8$ $U_8 = dD_8/dR_8 + N_8 dG_8/dR_8$ $X_8 = N_8 G_8 dF_8/dR_8$ and $Y_8 = -N_8 G_8 dM_8/dR_8$ (evaluated at the equilibrium). Note that $U_8 X_8$ and $Y_8$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_8$ as $(T_8)$ . (It follows that $f(T_8)$ 0 for all *T*, and $\int_0^\infty f(T_8)dt = 8.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_8^2 + U_8\lambda_8 - V_8\lambda_8 \int_0^\infty e^{-\tau_8\lambda_8} f(T_8) + X_8 + Y_8 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**ABSOLUTE STABILITY IN PREDATOR- PREY MODELS**

**APPLIED TO ACCENTUATION ATTRITION MODELS**

We apply model to each and every module and a concatenated model was to be presented which we procrastinate due to spatial restrictions.

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delay sand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - D(R) - NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR} > 0. \qquad (2)$$

The question is political-discursive. It's what Stengers and Latour are getting at with their concepts of diplomacy and cosmopolitics. They grant, first, that all entities exist and, second, that to say that someone's cherished idol (or whatever disputed entity they hold dear) is non-existent is a **'declaration of war'** – 'this means war,' as Stengers often says. They thus shunt onto-political discourse off of the terrain of knowledge/belief in the sense of existence/non-existence. Their basic claim seems to be that 'respect for otherness,' i.e. political pluralism, can only come from granting the entities that others hold dear an ontology, even if you don't 'believe' in them. You

are thus permitted to say 'I do not follow that god, he has no hold over me' but you are not permitted to say 'your god is an inane, infantile, non-existent fantasy, grow up.' And it's not just a question of politeness (although there's that too). The point is to grant others' idols and deities an existence – one needn't agree over what that existence entails, over what capacities that entity has or what obligations it impresses upon you as someone in its partial presence but to deny it existence entirely is to 'declare war' – to deny the possibility of civil discourse, of pluralistic co-existence. I believe that it was Richard Rorty who once quipped something like claims to reconcile realism and idealism always seem to end with a triumph of idealism. We don't, in fact, get realism through such approaches, but rather just get a pervasive anti-realism. I think this is also the problem with the "non-controversial pluralism" advocated by Stengers and Latour that Phillip defends here. Such pluralism is not realism but is, in fact, a thoroughgoing social constructivism. I think this is the central problem with Latour's argument in Irreductions (these days I regret having ever defended it). In rejecting both Enlightenment critique and what he calls "reduction" he wants to say something like "The Pentacostal really is filled with (e&eb) the Holy Spirit", that for the 19th working scientist heat really is a fluid and phlogiston really is what allows things to burn, and that for the Greek lightning really is an expression of Zeus's anger. Latour tells us that we aren't to reduce or explain away the entities posited by another group's "ontology" but are to develop explanations from within that ontology. It might not sound particularly sexy– and it certainly doesn't tell us what is worth thinking –but I can't help but believe that philosophy is the critical and reflective investigation of basic concepts that guide our investigation of the world about us, how we ought to live our lives, and what form of governance might be best. Compare two figures. A scientist might ask, what causes depression? We can very well imagine a philosopher turning around and asking the scientist, what is causality? The scientist presupposes a concept of causality in her investigations. She uses this concept in her inquiry. Now she might have a sophisticated concept of causality or she might never have thought much about causality at all, using it in the sort of colloquial and unreflective way that Plato decried when, for example, people

like Euthyphro talked about piety. A whole cascade of questions arise when we raise a question like "what is causality?" We can ask whether or not causality exists at all. We can ask how we distinguish between correlation and two events that merely accompany one another from genuine causation. This, for example, was Hume's question. But perhaps most importantly we can ask whether there is only one form of causality or many forms of causality. Is there only one-to-one causation; one cause and one effect? Is there many-to-one causation; or many events conspiring to produce an effect? Is there one-to-many causation; or one event producing a variety of different effects? We can even ask whether causality necessarily moves from past to present or whether there aren't forms of causality that move from future to past! **January 23, 2014 Pluralism and Realism:** multiple-worlds Over at Struggles Forever, Jeremy Trombley has an interesting post up on "the ontological turn" in anthropology or ethnography. I've been meaning to have a discussion with him about this as I think it's an issue many of us are struggling with. For example, the core project of The Democracy of Objects– a project which I think many have missed –is to somehow reconcile some version of social constructivism with a realist ontology capable of making room for ecology (which requires realist and materialist positions as there's a fact of the matter where global warming is concerned) as well as the role played by objective agencies in social assemblages such as technologies, infrastructure, features of geography, local climates, the growth cycles of plants and animals, waste, etc. Maybe we can try to organize some cross-blog event to discuss these issues. I certainly think they're close to the heart of Jeremy, Michael of Archive Fire, Arran James, and a host of others. As an aside, I'm beginning to realize how the different sites of the political I've been outlining– semiopolitics, thermopolitics, oikopolitics (political economy), geopolitics, eropolitics (the politics of sex and desire), biopolitics, and chronopolitics (and I'm sure there are other political sites!) –are drawing me away from traditional Marxism. Assuming that classical Marxism holds that economics or the conditions and relations of production are determinative of all other sites of the political, the various sites of the political that I've been outlining would lead to the conclusion that there is not one determinative

base of the political. This would not require committing Marx to flames, but rather of recognizing the phenomenon of overdetermination, or of a variety of different entangled sites of the political. But I digress. First, I find myself wondering what the ontological turn means in ethnography. Is it 1) the investigation of the different ontologies held/proposed by different cultures? E.g., the Aztecs believed that reality was structured in this way, while the Greeks in that way, and the ancient Chinese this way, etc? Or 2) is it an investigation of how real entities– independent of cultural beliefs –influence cultural formations? Or is it a combination of both? A position that I would favor. January 20, 2014 Model is applied to structured reality

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

**The evaporation spectrum of black holes from a local quantum gravity perspective Aurélien Barrau (Submitted on 21 May 2016**)

(15) Authors revisit the hypothesis of (e) a possible line structure in the Hawking evaporation spectrum of black holes, due to (e) non-perturbative quantum gravity effects, even arbitrarily far away from (e) the Planck mass.

(16) Authors show that this naive prediction might in fact hold in (eb)the specific context of loop quantum gravity, with (e&eb) a small departure from (e) the ideal case for (e) some low-spin transitions.

They also show that the effect is (=) neither washed out by the dynamics of (e) the process, nor by (e) existence of a mass spectrum up to (e) a given width, nor by the secondary component induced by (e) the decay of neutral pions emitted during (e&eb) the time-integrated evaporation. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Cite as: arXiv:1605.06613 [gr-qc] (or arXiv:1605.06613v1 [gr-qc] for this version

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t - T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\left(R(t)\right)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and Haimovici's model can be applied.

It's very easy to treat this as an abstract, academic debate: "Are you a realist or are you an anti-realist?", as if it were just a matter of what happens to be true. But it seems to me that this debate has, in reality, always been about politics. As I outlined in my last post, we have perpetually seen how appeals to the real and natural have been used in the name of oppressive power, inscribing both the exploitation of nature and the oppression of various people in the very fabric of being itself. Theistic theology and realist ontology have perpetually been used in the name of

what Deleuze called "State Philosophies" or philosophies that ontologize contingent orders of power and privilege (e.g., "the great chain of being" used to justify patriarchy, monarchy, serfdom, poverty, etc, and appeals to nature used to justify poverty and racial inequality (The Bell Curve), patriarchy (evolutionary psychology), heteronormativity, capitalism, etc). Because arrangements of power and inequality are always contingent in the sense that there's no marked difference in the capacities of peoples, power always looks for a transcendent supplement that would provide justification through ontological necessity. Antirealism– from the Greek atomists to present –became the radical and emancipatory gesture because it revealed the lie behind all of these forms of social organization or their inherit contingency or arbitrariness. Realism, by contrast, has all too often functioned as an apologetics for arbitrary power and social organizations. Here it's worth recalling what Foucault said about science: "…Even before we know to what extent something like Marxism or psychoanalysis is analogous to a scientific practice in its day-to-day operations, in its rules of construction, in the concepts it uses, we should be asking the question, asking ourselves about the aspiration to power that is inherent in the claim to being a science. The question or questions that have to be asked are: "What types of knowledge are you trying to disqualify when you say that you are a science? What speaking subject, what discursive subject, what subject of experience and knowledge are you trying to minorize when you say 'I speak this discourse, I am speaking a scientific discourse, and I am a scientist.' What theoretico-political vanguard are you trying to put on the throne in order to detach it from the massive, circulating, and discontinuous forms that knowledge can take?" (Society Must Be Defended, 10). All of these questions hold equally for claims to something being real. What is one trying to minorize when claiming something is real? What becomes privileged? What is excluded? It is these questions that have been at the heart of the realism debates, for as Spencer-Brown taught us, every distinction has a marked and unmarked space, draws attention to something to be included and pushes something into the unconscious or the domain of the invisible, hidden, or veiled. This is above all the case with evocations of the real. However, as I've tried to show antirealism

leads to its own problems. First, so long as we exclude real beings from our ontological inventory, we are unable to fully understand how power functions ("such and such a set of cultural formations have been thoroughly debunked, yet people still live as if they believed them"). Not only do we not fully understand the sources of the problems due to too much focus on the discursive and semiotic, but we deny ourselves valuable sites of political intervention at the level of infrastructure. Second, we are prevented from addressing things such as cultural racism, such as that found in Heidegger's privileging of the West and the Greeks and Germans in particular. We do a good job addressing biological and theological racism, heteronormativity, and sexism by showing how it is a cultural construction, but when faced with racism such as Heidegger's where he argues that there's something "unique" about the Greek event, it's language, and about the German language, or Badiou/Zizek's racism with respect to the "Pauline Event", we really have no response. Here someone like Jared Diamond or Fernand Braud el is needed to explain global-geographical inequalities. Third, the tools of the cultural turn really do not provide us with the means of thinking the ecological as a site of the political. For this reason, I've tried to formulate a third way– which might be called "constructivist realism" or "constructivist naturalism" –that retains the insights of the cultural turn, while also allowing a robust place for the material. I don't claim to be original in this. I think that many such as Manuel DeLanda, Deleuze and Guattari, Stacy Alaimo, Karen Barad, Jane Bennett, etc., are up to something similar. It's a vast project that requires the work of a multitude of voices, especially given the way in which the culturalism pervades humanities. **The Relational and the Non-Relational: Notes towards an Immanent and Pluralist Theory of Meaning Posted by larval subjects under uncategorized k-bigpic I have deleted or reformulated some sentences for which I beg pardon attributable to spatial constraints** What is correlation, what is causation? What is coordination? What is coincidence? These thoughts have bothered humans since long. I have been able to excavate some reliable material from Wikipedia which I would like to share. "All views are partial and contingent – that's the lesson of pluralism, and you say as much yourself when you say that we must be attentive to ways our own

knowledge might contain superstitions. But it's not discourse that determines whether one view works and another doesn't – it's a confrontation with **non-discursive (non-human)** agencies." I wonder if part of the issue here is that we understand pluralism differently. For me this doesn't sound like pluralism at all. Rather, it just sounds like our **epistemological condition**. We [hopefully] come to understand those portions of existence we question or investigate. The reason we come to investigate them is largely contingent. Finally, the accounts of these features of existence we give can be mistaken. All of this is perfectly consistent with a monism. I'm not sure why what he outlines above is the lesson of pluralism. It seems to me that every realist knows this. It seems to me that ontological pluralism is something quite different. Ontological pluralism is the thesis that there are many different worlds inhabited by many different entities. Thus, for example, you would have one world; say that of Lucretius, that's only inhabited by atoms and their combinations. You would have another world; say the world of the Mongolian shaman, that's inhabited by spirits that do all sorts of things. The ontological pluralist is saying that these spirits are, that they exist, that they're real. This is something quite different than merely saying that there are partial and contingent points of view on the world. It's this saying that these entities are rather than that some person or groups of people believe that they are that is the nub of the issue. Now, I think part of the issue here is that there's an ambiguity in the term ontology. Ontology can be one of two things. On the one hand, ontology is a group or persons set of beliefs as to what is. Here it's trivially true that there are pluralities of ontologies and the realist readily recognizes this. This is the whole reason there are debates over ontology. Mongolian shamans have their ontology, Europeans theirs, Christian fundamentalists theirs, materialists theirs, etc. When striving to understand and communicate with others it's vital to understand these ontologies because, as rhetoricians like Burke point out, our **beliefs about what is** are among the things that motivate our action. Despite having never seen bacteria I was my hands because I believe there are bacteria and viruses on door handles and whatnot and don't want to get sick. My belief about a particular thing existing is what motivates my action. And who knows, perhaps this belief is as

superstitious as the belief that the crops failed because God was displeased with my community. On the other hand, ontology is a theory **about what is.** It is making a claim that something exists. This is where the rubber hits the road. The ontological pluralist seems committed to the thesis that every group set of beliefs about what exists is sufficient for granting the existence of those entities. This is what I find objectionable in Latour. Obviously I'm not bothered by Latour's suggestion that we should take into account the role that nonhumans like speed bumps, rivers, microbes, etc., play in the form that social assemblages take. But this is not all that Latour claims. This is precisely where the philosopher might balk. Unlike the ethnographer, the philosopher is not interested in what people believe exists, but rather philosophers– at least of the realist variant –are trying to figure out what is. In other words, the realist philosopher begins with the premise that not all of these beliefs about what is are true. So for the philosopher, recognizing that Mongolian shaman's believe in the existence of shamans would only be the first step. The next step would consist in determining whether there's good reason for thinking such entities really do exist, i.e., whether there's good reason for believing these entities have mind and culture independent reality. Lest readers think that I'm just picking on the supernatural here, we can ask similar questions about strings, subatomic particles, galaxies, etc., etc., etc. I guess the question really comes down to what exactly we mean by ontological pluralism. When we talk about ontological pluralism are we defending the thesis that people have different theories of being? If so, then ontological pluralism is trivially true. If this is what is meant, then I certainly share Jeremy's view that it's valuable to understand the different world's people believe in. Certainly when I was practicing as an analyst I didn't get in ontological debates with my patients and it was necessary to understand their theory of being or their ontology to properly attend to them. Or, when we talk about ontological pluralism are we defending the thesis that all these ontological theories are true and refer to really existing entities? That's quite a different claim and is not one I would defend or endorse. Now someone might object that "in both The Democracy of Objects and Onto-Cartography you defend the thesis that there are multiple worlds." This is true. Because I hold that not everything

is related I'm led to the conclusion that there are diverse worlds. However, I also hold that however many worlds there might be, these worlds are nonetheless composed solely of material entities. Within the framework I propose I wouldn't suggest that there's one world where there are spirits and another world where there are souls and yet another where there are only material entities. My view is that there aren't spirits or souls in any of these worlds. One of the things I keep hearing in these discussions is that somehow the realist adopts a view from nowhere. I honestly don't understand this criticism. Investigation always occurs somewhere and requires all sorts of mediations involving technologies, experiments, etc. It's that labor of gathering evidence, conducting experiments, using technologies to observe the world, etc., that gradually gives us a body of data that allows us to say there's good reason to believe that such and such a thing exists and has these powers. Another charge seems to be that the realist refuses to recognize that their claims about the world are fallible. I find this charge particularly strange because it's precisely because the realist recognizes the difference between our theories of the world or what we say about the world and the world itself that fallibility is built into the core of his position. Realism doesn't mean one holds they have special access to the world, that they know all truths, or that they have the truth in hand, only that there are truths to be known and that we can be mistaken about things. Circling **Squares Philip** has a post responding to my quandaries about how to mesh realism and pluralism. He writes: Ontologically and metaphysically the idea of realist pluralism is no longer an issue. There are (appropriately) numerous variants but the basic idea that reality is itself pluralistic is well established. Captioned expatiation is about Everett worlds and string theory

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from

being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

The von Neumann formulation leads us to conclude that the lower realms of quantum reality form a virtually limitless array of potentials that are inherently incapable of realization without the observation of a knowing entity, which we identify with consciousness (Kafatos and Nadeau, 1990). We thus have a correspondence between quantum theory, the hierarchical structure implied by the Holographic Principle, and the higher orders of experience,

leading to the full expression of the Conscious Universe. David Bohm's theory of the implicate and explicate orders involves a holographic Principle that is fully consistent with the Holographic Principle discussed here (Bohm 1980; 1986; Bohm and Hiley, 1993). According to Bohm, there is an implicate order that represents the universal, holographic subtext of reality, and which unfolds in every moment to produce the explicate order that we all observe. Thomas Germinario (2004) has equated the implicate order with (eb) the unconscious process, and the explicate order with (eb) conscious process. He emphasizes the nature of dreams within (eb) the implicate order, and the importance of the dream work for maintaining a healthy mind through (e&eb) integration of the implicate order and our daily lives. Allan Combs and Mark Holland (1990) connected the implicate order or holomovement (Bohm, 1980) with (e&eb) Carl Jung is theory of synchronicity (Jung and Pauli, 1955), with (e&eb) the implicate order providing (eb) a holographic medium through which apparently disconnected individuals become connected. The principle of synchronicity, the instantaneous connection of people and events beyond (e) the senses, has been equated with the quantum-physical principle of non-locality (Combs and Holland, 1990; Germine, 1991), and has been proposed to be the fundamental mechanism of conscious process (Germine, 1991).Applying the Holographic Principle theory of mind, subjects may be connected synchronously in the manner of two individuals conversing through cellular phones, where the radio signals are transmitted through (e&eb) a distant satellite. The difference here is that the radio signals travel at the speed of light, and thus the there is some miniscule time lapse between the sending and receiving ends. The Holographic Principle implies (eb) a much more distant and instantaneous communication pathway, with (e&eb) the analogue of the satellite being the holographic boundary of the Universe, as well as a much richer communication, potentially involving a transmission of thoughts and feelings. Our ego-consciousness seems to mask (e) the universal relatedness implied by (e) the Holographic Principle, and it is perhaps only through transcendence of (e) the ego-consciousness that the higher orders of experience can become conscious. In the phenomenon of synchronicity, there seems to be a meaningful connection between (e&eb) individuals that

breaks through the barrier of ego-consciousness. Such a connection is reported by many individuals in the course of dreams, when (e) the ego-consciousness has been suspended, at times informing (e&eb) the dreamer of something that has happened in the life of a meaningfully-connected individual. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Impassioned-intense Connectionist neuroscience, in an effort to attain relevance to clinical practice, has recently given us a theory of the relationship of the unconscious to (e&eb) consciousness which is (=) dualistic (Viamontes and Beitman, 2007), an ad hoc theory of conclusions based on (e) assumptions, the most erroneous of which are that representations of (e) information are information, and that such representations elaborate themselves through (e&eb) a process of representations of representations in neuronal circuits (Bechtel and Abrahamsen,1991). As theory, the connectionist models are lacking (e) heuristic value and **are (=) inimical to the progress** of (e) basic and clinical neuroscience. This is not to denigrate or minimize the importance of neuronal connections, but these connections are rather like the wires in (eb) a radio that plays a symphony. The radio, in and of itself, is (=) incapable of playing (e&eb) the symphony. The radio doesn't do not write the symphony, nor does it broadcast the Passionate-raging symphony over the airways. Yet, if we loosen one wire, the symphony is (=) not manifested through the radio. The brain is (=) not subject to (e) the same vulnerability, in that it has (e) parallel and redundant networks of processing, and, it is in this sense that the connectivity of the brain becomes (eb) important. Connectionism, to the extent that it accepts (e) consciousness and the (e&eb) unconscious at all, attributes them to (e) separate groups of "circuits" in the brain, with (e&eb) consciousness-processing circuits having (e) a limited but detailed ability to analyze (e&eb) information, and unconscious-processing processing more information in less detail (Viamontes and Beitman, 2007). There is an element of truth in this view in that information that has importance and meaning resonates within the brain at higher levels of experience and at higher levels of holographic recursion. This resonance involves what we know, what we feel, and what we are capable and willing to consciously realize. The Universe, as William Blake noted,

can be seen in a grain of sand. Indeed, the grain of sand requires (e) the Universe to exist. The human organism requires (e) the Self to exist. The Self is not a simple amalgam of (e) sensory experiences, as suggested in the connectionist model (Blinder, 2007). This amalgam view of the Self can be damaging when applied to (e&eb) psychotherapy, as the connectionists suggest (Blinder, 2007). The Self is ours forever and ever, from eternity to eternity. As Milton suggested, the mind is (=) its own place. Our deepest purpose is for our minds to resonate with (e&eb) the supra-conscious levels of experience, much as it is (=) purpose of the radio to play the symphony. If the symphony is not to our liking, (eb) we can turn the radio off. So, perhaps, the transpersonal consciousness can turn (e&eb) our own individual and collective radios off, or perhaps turn the volume down, if (e) we are not in harmony with the **Universal Mind**. Physics has given us many enigmas. Quantum theory itself is an enigma. But some of the enigmas seem to involve (e&eb) the limitation of (e) our own abilities of discernment. The thought experiment of (e) "Schrödinger's Cat," as it is often interpreted (e.g. Gribbin, 1984), assumes (eb) that the **cat is not conscious, and is (=) incapable of observation in the context of (e) collapse of (e) the wave function.** This is due, in large part, to the false dichotomy of (e) information and (e&eb) experience. Our observation of other animals indicates, to those that are attuned to (e) the animal mind, who can "feel" the animal mind, that, at (eb) the very least, all mammals are (=) conscious. It would seem that the great neuropsychiatrist, Stanley Cobb (1948), may have been correct in his attribution of (e) levels of consciousness to a wide variety of organisms: "...lower animals with (e&eb) no cerebrum appear to be conscious...even plants such as the sunflower that turns towards strong light may have (e) a vague awareness or warmth and comfort. There are many degrees of consciousness and it is my contention that it is integrated at (e&eb) many levels like other important functions of (e) the central nervous system."With respect to the theory of mind, recent mainstream thinking remains (eb) classical and mechanistic. Such views of mind are advocated by our most in prominent neuroscientists (e.g. Changeux, 1985; Gazzaniga, 1985) as neuroscience research attempts to explain all mental function on the basis of (e) **pure mechanism and the localized function of**

**specific areas of the brain.** Francis Crick, the distinguished Nobel Laureate, had a very influential second career as a neuroscientist. He concluded that we are basically soulless creatures, that all mental processes could be reduced to (e&eb) identifiable, neural correlates in (eb) the brain, and that, in fact, we are (=) nothing but a collection of neurons (Crick, 1995). Crick even went so far as to say that what we call the soul is (=) a group of neurons located in the prefrontal lobes of the brain. Another Nobel Laureate, Gerald Edelman (1987), has written a critically acclaimed book on "neural Darwinism," which purports that the brain is "circuitry" is constructed through a developmental process of "survival of the fittest" neurons and neuronal connections. At the same time the total purposelessness and lack of direction in evolution is expounded, again to critical acclaim, by Richard Dawkins (1986). The idea that there essentially no self or soul imbued with agency and self determination was expounded by Gilbert Ryle (1949/2002) with his parody of the "ghost in the machine," and this work continues to be influential among many scholars of the mind. Philosophers such as Patricia Churchland (1986), who coined the term "neurophilosophy," take an "eliminative materialist" view of (e) belief, free will, and consciousness. Other philosophers, such as Daniel Dennett (1991) would make (eb) experience "epiphenomenal," a by-product of (e) brain processes with (e&eb) no effect whatsoever on (eb) the function of mind. However, unlike neuroscience, which is still in its infancy, the philosophy of mind has had a long history, including such notable figures as Plato, Aristotle, Descartes, Kant, Locke, Hume, Kierkegaard, and James, to name just a few. Perhaps, as noted by Ervin Laszlo (1974): "in today's world, most of the traditional functions of (e) cognitive synthesis have atrophied (e) and are ignored and neglected." Hopefully, what we have presented here will be a step in the right direction for what Laszlo calls a "conceptual synthesis" to help fill the need for (e) meaningful engagement is such a world. The fundamental elements of such a synthesis have been described by (e) Laszlo as follows: Conceptual synthesis performs (e&eb) at least five basic functions in the guidance of affairs. They are the mystical, the cosmological, the sociological, the pedagogical or psychological, and the editorial functions. The mystical function inspires (eb)

in man a sense of mystery and profound meaning related to (e&eb) the universe and of himself in it. The cosmological function forms (eb) images of the universe in accord with (e&eb) local knowledge and experience, enabling (eb) men to describe and identify (e&eb) the structure of (e) the universe and the forces of nature. The sociological function validates supports and enforces (eb) social order, representing (eb) it in accord with (e&eb) the nature of the universe, or as the natural or right form of (e) social organization. The pedagogical or psychological function guides (eb) individuals through stages of life, teaching ways of understanding themselves and others and presenting (eb) desirable responses to life is challenges and trials. Finally, the editorial function of (e) conceptual is to define (eb)some aspects of reality as important and credible and hence to be attended to, and other aspects unworthy of serious attention. **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science** Kant sums up this inversion and its spirit early in the Critique of Pure Reason: Up to now it has been assumed that all our cognition must conform to the objects; but all attempts to find something about them a priori through concepts that would extend our cognition have, on this presupposition, come to nothing. Hence let us once try whether we do not get farther with the problems of metaphysics by assuming that the objects must conform to our cognition, which would agree better with the requested possibility of an a priori cognition of them, which is to establish something about objects before they are given to us. This would be just like the first thoughts of Copernicus, who, when he did not make good progress in the explanation of the celestial motions if he assumed that the entire celestial host revolves around the observer, tried to see if he might not have greater success if he made the observer revolve and left the stars at rest. (B xvi) Onticology, like all variations of object-oriented ontology, is realist in its orientation. In defending a realist ontology onticology holds that the vast majorities of objects, actants, beings, or entities are independent of humans and are what they are regardless of whether any humans regard them or register them. In short, onticology rejects any anthropomorphic, idealist, or anti-realist thesis to the effect that to be is to be the correlate of mind, spirit, the body, the human, and language or otherwise. While it is certainly the case that knowledge is necessarily dependent on the object to which it relates, the converse does not hold true. Objects are not dependent on being known, regarded, perceived, or spoken

about. As such, and to put it in Aristotlean terms, knowledge is an accident of objects, not objects an accident of knowledge. As Althusser so nicely puts it, "[n]o doubt there is a relation between thought-about-the-real and thisreal, but it is a relation of knowledge, a relation of adequacy or inadequacy of knowledge, not a real relation, meaning by this a relation inscribed in that real of which the thought is the (adequate or inadequate) knowledge" (Reading Capital, 96). Althusser goes on to remark that "[t]he distinction between a relation of knowledge and a relation of the real is a fundamental one: if we did not respect it we should fall irreversibly into either speculative or empiricist idealism" (ibid.). Onticology categorically endorses Althusser's verdict. It is a fundamental necessity to distinguish between those relations that belong to the object and those that belong to knowledge. Contemporary philosophy continuously confuses these two very different sorts of relations. Naturally the question arises of how it is possible to surmount our relation to the object so as to determine whether objects themselves possess the properties we encounter in relating to objects. In other words, given that we can only ever relate to the object in relating to the object how is it possible to surmount this relation to get at the being of the object itself? Much more will have to be said about this later– and the answers will be surprising with respect to standard prejudices about realism –however, for the moment it can be said that onticology takes its epistemological inspiration from the transcendental realism of Roy Bhaskar. Among other things, Bhaskar sought to provide a transcendental grounding for the sciences. Insofar as onticology defends the thesis that the field of being is much more vast than the field of objects investigated by the natural sciences, it parts way with the thesis that the domain of being is exhausted by the domain of natural objects. However, the general form of Bhaskar's argument holds for our realist purposes. A transcendental argument seeks to elucidate the conditions under which certain acknowledged practices and forms of cognition are possible. Kant, for example, asked what must be the case for mathematical judgments to be possible. How is it both that we are able to extend our knowledge, as if by magic, through mathematical judgments and, more significantly, that these judgments are able to provide genuine knowledge of the world despite the fact that these

forms of reasoning are not based on experience? Part of Kant's argument consisted in claiming that mind imposes the forms of space and time on the data of experience. In other words, space and time are not attributes of being itself but rather of the mind that regards being. Insofar, Kant argues, as mathematics is ultimately a rumination on the nature of space and time taken in their most abstract form and insofar as the mind imposes space and time on the manifold of sensation, it thus follows that a priori judgments about the nature of spatio-temporal relations are possible that anticipate the structure of actual-space times without directly experiencing these space-times. Why? Because any manifold of sensation must necessarily be structured by these forms imposed by intuition. **Onticology–A Manifesto for Object-Oriented Ontology Part I Posted by larval subjects under Object-Oriented Philosophy My Question Is This: You Say, "Granted There Are Deeper And Deeper And More Primitive Forms Of "Knowledge" Or Abstraction Into Mechanism, But I Wouldn't Say That The Atomic Level Brings A New Epistemic World Into View." But Doesn't Any Kind Of Knowledge, However Primitive, Imply Bringing Some Kind Of Epistemic World Into View Or Proto-View?** Joel: David, Yes Exactly. My Question Is, Just How Far Down The Complexity Gradient Can The Meaning Of 'Knowledge,' Or 'Semiotics, Or Specifically Representation (Maya) Be Stretched Without Breaking? I Can Comfortably Stretch It To The Genetic Level, As The Proto-Epistemic, Where Genetic Codes 'Represent' Phenotypic Structures Or Modes, And The Whole Of Evolution Learns From Its Mistakes And Successes As That Knowledge And History Is Encoded Into And As The Recapitulation Of The Embryogenesis Of The Organism. But An Atom, In My Model, While Certainly Reacting To Its Environment From Its Own Attractor, Essence, And 'Interiority' (Prehension) Isn't Making Representations Of It To Do So. But The Key To The Power And 'Essence' Of Representation (And Simultaneously Its Limitation) Is The Capacity Of Abstraction And Choice. Evolution Does This At A Rudimentary Level (And Anthropomorphizing A Bit) With Its Sending Into The World Its 'Random' Variations, Organismic Ideas, In A Sense. Its 'Mind' Is Literally On The Outside As Our Living World And Biosphere. A 'Choice' Is Made In The Biospheric Interiority,

And Knowledge For The Evolutionary/Embryogenetic Trajectory Is Gained By The "Differential Reproductive Success" Of These Individuals Lives Through Time, As The Increasing Intelligence, Representation, And Capacity For Choice Is Directly Injected Into The Flow Through Sexual Selection. But Where Is The Encoded/Abstracted (Enfolded) Representation Of Options And A Choice Among Them Made At The Atomic Level? I Am Working With The Most Cutting-Edge And Coherent Models For The Atom (Including My Own) That I Can Find Among The Heterodoxy, And I See No Evidence Of Representation And Choice At This Level, Or Rather No Real Way To Further Stretch The Meaning Of The Terms To This Level. Sure There Is Uncertainty From Infinite Difference And Immanent Causation (See Bohm On Infinite Causation), But Without Representation There Really Is No Choice (Although Crudely We Could Call A 'Bifurcation' A 'Choice' By The System As A Whole). So I Am Simply Saying That Representation At The Atomic Level Is Entirely Enfolded And These Are Really Just Ontic Level Phenomena. The Atomic Level, In My View, Is Pre- Not Yet Proto-Epistemic, Which I Place at the Genetic Level Where We Find the First Codes And (e&eb) Primitive Representation. This Just Means That There Is A Real Distinction Between Ontic And Epistemic (And A Gradient Between Them) And We Can't Collapse The Two Into Just The Ontic-Epistemic. And Yet, With The Distinction Intact, And The Gradient To Explore, They Remain Nondual, Univocal, ONE Or Ontic. What I Would Say Exists At All Levels (My Own Self-Aware Myth About The Given) Is What I Call The "Symbiogenesis Of Subject And Object". This Is The Spinozan Essence Of Dynamic Stability And Growth, And Leibnizian Prehension. Also Called The "Nucleation Of Observability" In Spinbitz. It Is From This Nucleus Of The Observer (Fuller) In A "Point-Free Geometry" (Whitehead) Of Pointless "Points Of View" As Conditions Of Boundary That Wilber's "All Is Perspective" Finds Its Ontic Grounding. And This Symbiogenesis Is Also A Rudimentary Or Deep-Level Native Or Primitive Intelligence, Evolution At The Involution Of The Subject-Object Interface. It Is An Exploration Of Creative Learning And Direct Immanent Awareness, Unmediated (Enfolded) By Representational Forms (Unfolded). It Just Is, In Waves (Cosmic Vertebrae)

Before It Abstracts, Represents And Incarnates The Flesh Of What Could Be. Matter Is Enfolded Seed Of The Flower Of Its Abstraction. Maya Is This Depth And Recursion Of Brahma Into And Through The Boundary Conditions Of Abstract Relation. And Brahma Is The Span Of The Emergent Rigor Of The Cosmic Ergodic Spine That Opens The Recursions Into The Field Of Infinite Difference. In Immanent-Transcendent Waves, Literally And Empirically A Limit-Cycle In Ergodicity, It Enfolds And Unfolds From Anatom To Anatomy, From Infinite Intelligence Or Omnirelational Awareness Of Self, Through The Representational Depths Of Self As Self-Consciousness. Omni-Evolution, Trans-Dynamics, And The Anatom Posted On February 21, 2013A Discussion With Tom Huston, Part II The Following Is Part 2 Of An Excerpt From A Discussion Initiated By Tom Huston On His Facebook Page, Edited And Embellished For Clarity. Find Part One, Here. In This Part We Dig Into The Nature Of The "Anatom" And The "Ergodic Spine" Of An Electro-Fractal Cosmos, And Into The Trans-Dynamic Integration Between Being And Becoming, And The X-Interface−The Crossroads Of The Ontic-Epistemic (Brahma And Maya) And The Subject-Object Polarities. Univocal_Dynamics_V1 Tom Huston Is A Founding Member Of Integral Institute And A Former Editor Of Enlighten Next Magazine. His Webpage Is Tomhuston.Com. David Marshall Is A Writer And Editor Living In Chicago. No meaning, value, or normativity to be found in nature, that there is nothing **natural beings ought to be,** but that, rather, these judgments arise from us (See last post of larval subjects) which has led some to raise valuable questions about the coherence of these claims. The problem arises when the following three propositions are taken together: 1. there is nothing outside of nature.2. Beings have no intrinsic meaning, purpose, or value (in and of themselves, there's nothing they ought to be). 3. Value judgments about what beings and being ought to be arise from us and beings like us (bonobo apes, dolphins, institutions, birds of paradise, etc). The problem arises between proposition 1 and 3. How can it both be true that there is nothing outside of being and those normative judgments belong to us and other beings capable of making normative judgments, not nature? The problem arises from restricting these judgments to humans and beings capable of

making these judgments. In making such a claim it seems as if we're saying that there's something outside of nature, something that is beyond nature, thereby violating the first thesis and potentially reintroducing the nature/culture distinction. By contrast, if we say that normative judgments are the special domain of those living beings with the proper degree of sentience to make such judgments, then we seem to reintroduce the nature/culture distinction and fall back into the sorts of problems that I outlined in my last post (and that are so nicely critiqued by thinkers such as Latour). Is there a way out of this? I don't know. Great_Chain_of_Being_2_ (lighter) ASIDE: It seems to me that this is really what the debate between realism and anti-realism, realism and socio-linguistic constructivism surrounding the new materialisms and speculative realists has really been all about.

Opposite limit of absolute stability with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

If a predator or a debit types takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \qquad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9}+\frac{N_9 dG_9}{dR_9}-\frac{dB_9}{dR_9}>\frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9}+\frac{dM_9}{dR_9}\right], \qquad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usance bills bin according with terms and conditions of the letter of credit.

**ABSOLUTE STABILITY**

$$L_9=\frac{N_9 G_9}{2M_9}\frac{d(9nF_9)}{dR_9}(\geq 0) \qquad (9a)$$

$$I_9=-\frac{N_9 G_9}{2M_9}\frac{d(9nM_9}{dR_9}(\geq 0) \qquad (9b)$$

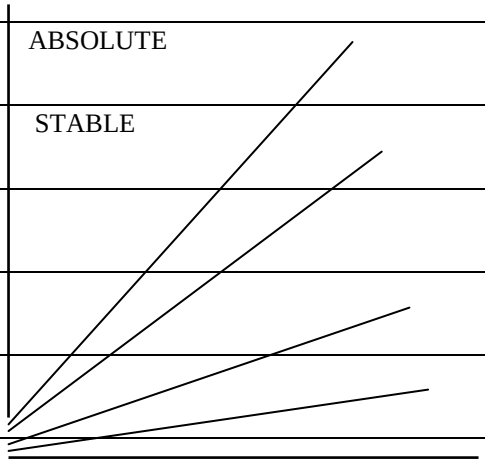
It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9>L_9-I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$$P_9>L_9-I_9, \quad \text{when} \quad P_9^2/4>I \qquad (8a)$$

$$P_9>L_9-I_9+\frac{(I_9-P_9^2}{(L_9-I)} \quad \text{when} \quad P_9^2/4<I. \qquad (8b)$$

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model[$c.f.condition(2)$]. The delayed term, L, acts directly to reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,




DELAYED PREDATOR EFFECT

FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect ($I = (NGdM/dR)/2M^2$).

Reffer couffesy Nunney

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability

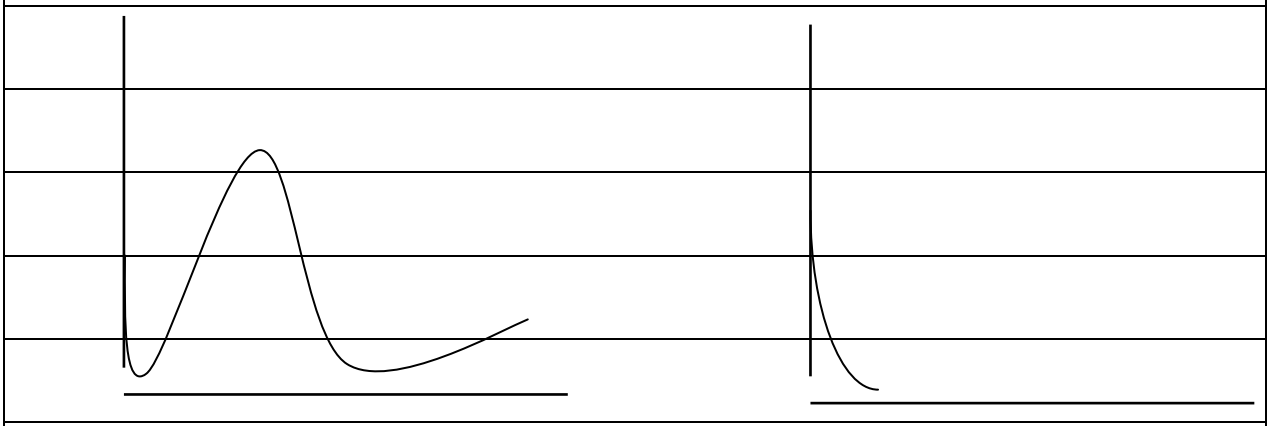
 

is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the



TIME DELAY

Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

Indeed, Nunny's model can be applied to some of the consciousness and philosophical problems. Predator prey analysis has ripped open the veneer of the fundamental concept of universe namely Debit-Credit world.

The noise of ordinary consciousness and (e&eb) beliefs drowns out (e) the signal. Unconsciousness is the background of (e) our ordinary awareness. Our organism is very much at (eb) the center of such effects. The organismic source is (=) our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is a meme-Scape. Pre-conceived concepts vie with structures, concepts with images. Like scientists who ignore assumed truths, we leapfrog over our beliefs and personality deficits, claiming idiosyncratic imagination is literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is both. We inherit the structures that make our experience what it is.

But the structure itself is “empty,” and each human culture “fills” it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us

feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through 're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns, nomenclature, and internal coherence of meaning and representation. In a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist

simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** The lower domain of experience is what we do not directly observe, the quantum realm, while the higher level is what we ordinarily observe, the classical realm. Based purely on experiment, the formulation of quantum theory initially placed the great divide between the quantum and classical domains between the measuring apparatus and the particle (Herbert, 1985). However, in the early 1930s, John von Neumann, in his rigorous mathematical treatise on quantum mechanics, found no support for such a division, compelling him to conclude that the wave for was collapsed by consciousness. Von Neumann's treatise has been called the "quantum bible" and "the most influential book on quantum theory ever written" (Herbert, 1985).

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t-T_9)F_9\big(R_9(t-T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t-T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ and the effective prey regulation $(dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9)$ must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting $dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

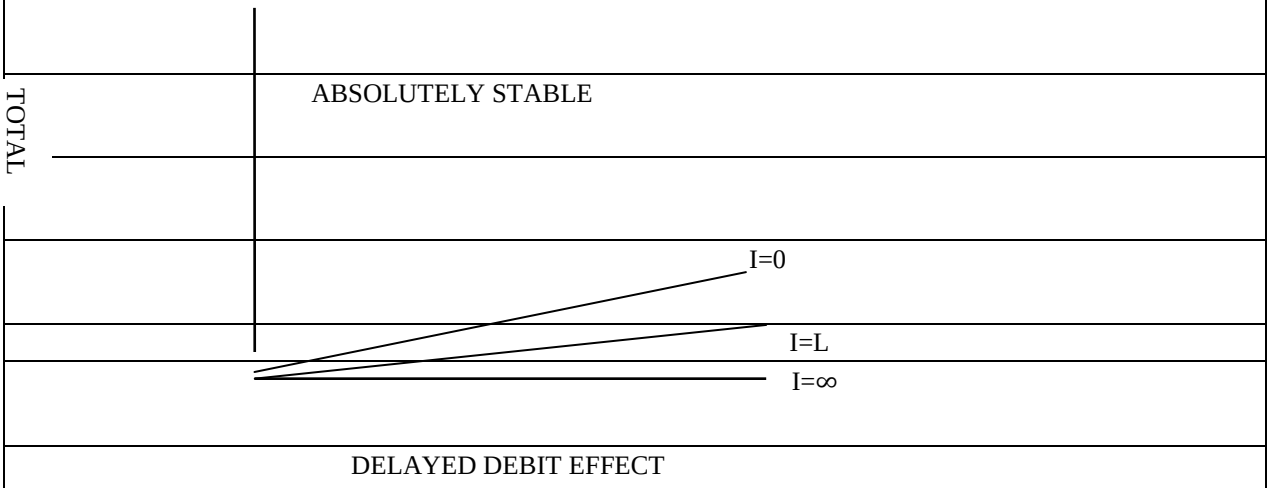



FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.

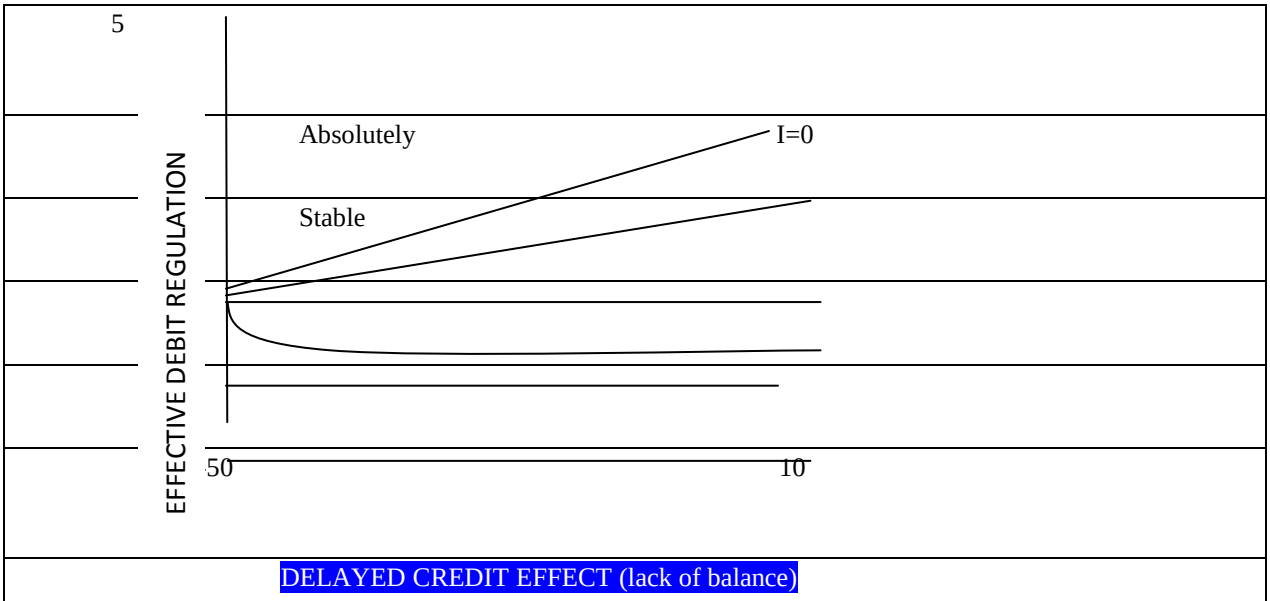



Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the

constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^\infty \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9\left(U_9 - V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9\right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9\int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$ then given the constraints of $f(T_9)$, $(A1.5)$ *can never be satisfied* Consequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

**ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDERATION**

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 - V_9M_9 -]e^{-T_9\lambda_9} = 0$$

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9 Z_9 \sin(Z_9 T_9) - (A_9 M_9 + Y_9) + (A_9 M_9 + X_9)\cos(Z_9 T_9) = 0$$

$$(A_9 M_9) Z_9 - M_9 Z_9 \cos(Z_9 T_9) + (A_9 M_9 - X_9)\sin(Z_9 T_9) = 0$$

Where $A_9 = U_9 - V_9$ Using (A2.3), substitute for $Z_9^2$ and $Z_9$ in (A2.2):

$W_9^2(X_9^2 - 3M_9 A_9 X_9 - M_9^2(X_9 - Y_9) + 2M_9^3 A_9 + 2A_9^2 M_9^2]$

$$+W\left[MA^{3}-A^{2}\mathrm{X}+2\mathrm{MA}(2\mathrm{X}+\mathrm{Y})-2X^{2}\right]+A^{2}(\mathrm{X}+\mathrm{Y})=0$$

*where* $W_9$=1-cos ($\mathrm{Z}_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (\mathrm{X}_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ \mathrm{X}_9 - Y_9/2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(\mathrm{W}_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \le W_9 \le 2$ *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (\mathrm{X}_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\,\mathrm{M}_9)^2\,(2A_9\,\mathrm{M}_9 - (\mathrm{X}_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9M_9^2 - 3M_9X_9\right) + W_9\left(3M_9A_9^2 - 2A_9X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9 + Y_9\right)$$

$=2M_9^2W_9^2(M_9 + A_9) + M_9A_9^2W_9 + \left(2A_9M_9 - X_9 + Y_9\right)(A_9W_9 + M_9W_9^2$

$+ (2- W_9)((2M_9X_9 + M_9Y_9)W_9 + A_9(X_9 + Y_9))$,

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range 0≤ $W_9$ ≤2.Defining

$Q_9(W_9, A_9) = a_9W_9^2 + c_9$ *where* $a_9$ $b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

**Harald Atmanspacher, Hans Primas** show that such an epistemic perspective can be consistently complemented by Einstein's ontically oriented position. **Time, Quantum and Information 2003, pp 301-321 Epistemic and Ontic Quantum Realities Harald Atmanspacher, Hans Primas Robert W. Spekkens** present a toy theory that is based on a simple principle: the number of questions about the physical state of a system that are answered must always be equal to the number that are unanswered in a state of maximal knowledge. Many quantum phenomena are found to have analogues within this toy theory. These include the noncommutativity of measurements, interference, the multiplicity of convex decompositions of a mixed state, the impossibility of discriminating nonorthogonal states, the impossibility of a universal state

inverter, the distinction between bipartite and tripartite entanglement, the monogamy of pure entanglement, no cloning, no broadcasting, remote steering, teleportation, entanglement swapping, dense coding, mutually unbiased bases, and many others. The diversity and quality of these analogies is taken as evidence for the view that quantum states are states of **incomplete knowledge** rather than states of reality. A consideration of the phenomena that the toy theory **fails to reproduce, notably, violations of Bell inequalities** and the existence of a Kochen-Specker theorem, provides clues for how to proceed with this research program.DOI: http://dx.doi.org/10.1103/PhysRevA.75.032110 **Evidence for the epistemic view of quantum states A toy theory Phys Rev A 75, 032110 – Published 19 March 2007 Robert W. Spekkens In such specialized and autonomous areas, it is to the eminent and erudite authors to interpret the contradistinctions and contradictions' in the dove tailing explaining. This I feel would be yeoman service in the creation of Nature's general ledger. M. S. Leifer and O. J. E. Maroney** examine the relationship between quantum contextuality (in both the standard Kochen-Specker sense and in the generalized sense proposed by Spekkens) and models of quantum theory in which the quantum state is maximally epistemic. We find that preparation noncontextual models must be maximally epistemic, and these in turn must be Kochen-Specker noncontextual. This implies that the Kochen-Specker theorem is sufficient to establish both the impossibility of maximally epistemic models and the impossibility of preparation noncontextual models. The implication from preparation noncontextual to maximally epistemic then also yields a proof of Bell's theorem from an Einstein-Podolsky-Rosen-like argument. DOI: http://dx.doi.org/10.1103/PhysRevLett.110.120401 Received 25 August 2012Published 20 March 2013 © 2013 American Physical Society **Maximally Epistemic Interpretations of the Quantum State and Contextuality Phys. Rev. Lett 110, 120401 – Published 20 March 2013 M. S. Leifer and O. J. E. Maroney.** Epistemological, onitic and ontological aspects particularly of quantum mechanics have invited

profound thinkers to contribute their mite to the growing literature on the subject, which includes the contributions from the subterranean realm and ceratoid dualism of philosophy like that Chaos theory of Deleuze, Derrida , Kierkegaard and others. **Andrei Khrennikov** shows that the Dirac-von Neumann formalism for quantum mechanics can be obtained as an approximation of classical statistical field theory. This approximation is based on the Taylor expansion (up to terms of the second order) of classical physical variables – maps f: Ω → R, where Ω is the infinite-dimensional Hilbert space. The space of classical statistical states consists of Gaussian measures ρ on Ω having zero mean value and dispersion σ2 (ρ) ≈ h. This viewpoint to the conventional quantum formalism gives the possibility to create generalized quantum formalisms based on expansions of classical physical variables in the Taylor series up to terms of nth order and considering statistical states ρ having dispersion σ2(ρ) = h n (for n = 2 we obtain the conventional quantum formalism). **December 2005, Volume 18, Issue 7, pp 637-650 Date: 28 Nov 2005 Generalizations of Quantum Mechanics Induced by Classical Statistical Field Theory Andrei Khrennikov** It is widely accepted that consciousness or, in other words, mental activity is in some way correlated to the behavior of the brain or, in other words, material brain activity. Since quantum theory is the most fundamental theory of matter that is currently available, it is a legitimate question to ask whether quantum theory can help us to understand consciousness. Several approaches answering this question affirmatively, proposed in recent decades, will be surveyed. It will be pointed out that they make different epistemological assumptions, refer to different neurophysiological levels of description, and adopt quantum theory in different ways. For each of the approaches discussed, these imply both problematic and promising features which will be indicated. **Discrete Dynamics in Nature and Society Volume 2004 (2004), Issue 1, Pages 51-73 http://dx.doi.org/10.1155/S102602260440106X Quantum theory and consciousness: an overview with selected examples Harald Atmanspacher Copyright © 2004 Hindawi Publishing**

**Corporation. This is an open access article distributed under the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original work is properly cited.** Another thought provoking and investigatory article is by Peter G. Lewis et al. Perhaps the quantum state represents information about reality, and not reality directly. Wave function collapse is then possibly no more mysterious than a Bayesian update of a probability distribution given new data. **Peter G. Lewis et al** consider models for quantum systems with measurement outcomes determined by an underlying physical state of the system but where several quantum states are consistent with a single underlying state—i.e., probability distributions for distinct quantum states overlap. Significantly, they demonstrate by example that additional assumptions are always necessary to rule out such a model. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.150404 **Distinct Quantum States Can Be Compatible with a Single State of Reality Phys. Rev. Lett. 109, 150404 – Published 9 October 2012 Peter G. Lewis et al** Communication complexity of a quantum channel is the minimal amount of classical communication required for classically simulating a process of state preparation, transmission through the channel and subsequent measurement. It establishes and perpetuates a limit on the power of quantum communication in terms of classical resources. **Alberto Montina** shows that classical simulations employing a finite amount of communication can be derived from a special class of hidden variable theories where quantum states represent statistical knowledge about the classical state and not an element of reality. This special class has attracted strong interest very recently. The communication cost of each derived simulation is given by the mutual information between the quantum state and the classical state of the parent hidden variable theory. Finally, find that the communication complexity for single qubits is smaller than 1.28 bits. The previous known upper bound was 1.85 bits. DOI: http://dx.doi.org/10.1103/PhysRevLett.109.110501 **Epistemic View of Quantum States and**

**Communication Complexity of Quantum Channels Phys. Rev. Lett 109, 110501 – Published 12 September 2012 Alberto Montina for models see elsewhere .** Stability analysis of what? Quantum state representing reality or our knowledge of reality? Does the quantum state represent reality or our knowledge of reality? In making this distinction precise, **Nicholas Harrigan, Robert W. Spekkens** are led to a novel classification of hidden variable models of quantum theory. They show that representatives of each class can be found among existing constructions for two-dimensional Hilbert spaces. Approach also provides a fruitful new perspective on arguments for the nonlocality and incompleteness of quantum theory. Specifically, **Nicholas Harrigan, Robert W. Spekkens** show that for models wherein the quantum state has the status of something real, the failure of locality can be established through an argument considerably more straightforward than Bell's theorem. The historical significance of this result becomes evident when one recognizes that the same reasoning is present in Einstein's preferred argument for incompleteness, which dates back to 1935. This fact suggests that Einstein was seeking not just any completion of quantum theory, but one wherein quantum states are solely representative of our knowledge. Hypothesis is supported by an analysis of Einstein's attempts to clarify his views on quantum theory and the circumstance of his otherwise puzzling abandonment of an even simpler argument for incompleteness from 1927. **February 2010, Volume 40, Issue 2, pp 125-157 Date: 09 Jan 2010 Einstein, Incompleteness, and the Epistemic View of Quantum States Nicholas Harrigan, Robert W. Spekkens** statistical model of the probabilistic description of physical reality has been studied by many authors. **Andrei Khrennikov posits** a contextual statistical model of the probabilistic description of physical reality. Here contexts (complexes of physical conditions) are considered as basic elements of reality. There is discussed the relation with QM. He proposes a realistic analogue of Bohr's principle of complementarity. In the opposite of the Bohr's principle, **Andrei Khrennikov** principle has no direct relation with mutual exclusivity for

observables. To distinguish his principle from the Bohr's principle and to give better characterization, he changes the terminology and speaks about supplementarity, instead of complementarity. Supplementarity is based on the interference of probabilities. It has quantitative expression trough a coefficient which can be easily calculated from experimental statistical data. There is need not appeal to the Hilbert space formalism and noncommutativity of operators representing observables. Moreover, in our model there exists a pair of supplementary observables which cannot be represented in the complex Hilbert space. There are discussed applications of the principle of supplementarity outside quantum physics. **October 2005, Volume 35, Issue 10, pp 1655-1693 The Principle of Supplementarity: A Contextual Probabilistic Viewpoint to Complementarity, the Interference of Probabilities and Incompatibility of Variables in Quantum Mechanics Andrei Khrennikov Foundations of Physics** Description of the vacuum in Yang-Mills theory and arrive at a physical interpretation of the pseudoparticle solution and the attendant violation of symmetries has been done by **R. Jackiw and C. Rebbi.** It is here vaccuum stability forms the bastion, pillar, post and stylobate of the mechanisms of yang mills Theory and action of a particle. The existence of topologically inequivalent classical gauge fields gives rise to a family of quantum mechanical vacua, parametrized by a CP-nonconserving angle. The requirement of **vacuum stability** against gauge transformations renders the vacua chirally noninvariant. DOI: http://dx.doi.org/10.1103/PhysRevLett.37.172 **Vacuum Periodicity in a Yang-Mills Quantum Theory Phys. Rev. Lett 37, 172 – Published 19 July 1976 R. Jackiw and C. Rebbi.** Three-dimensional Yang-Mills and gravity theories augmented by gauge-invariant mass terms are analyzed. These topologically nontrivial additions profoundly alter the particle content of the models and lead to (eb) quantization of a dimensionless mass-coupling-constant ratio. The vector field excitations become (=) massive, with (e&eb) spin 1 (rather than massless with spin 0), and the mass provides (eb) an infrared cutoff. The gravitation

acquires (e) mass, mediates (e&eb) finite-range interactions, and has (e) spin 2 (rather than being absent altogether); although its mass term is of (e)third derivative order, there are no ghosts or acausalities. DOI: http://dx.doi.org/10.1103/PhysRevLett.48.975 © 1982 The American Physical Society **Three-Dimensional Massive Gauge Theories Phys. Rev. Lett 48, 975 – Published 12 April 1982 S. Deser et al** All scattering amplitudes in (eb) the maximally supersymmetric N=4 super-Yang–Mills theory possess (e) a new, dual superconformal symmetry which extends the previously discovered dual conformal symmetry of MHV amplitudes. To reveal this property authors formulate the scattering amplitudes as (=) functions on (e&eb) the appropriate dual superspace. Rewritten in this form, all tree-level MHV and next-to-MHV amplitudes exhibit manifest dual superconformal symmetry. Authors propose a new, compact and Lorentz covariant formula for (e) the tree-level NMHV amplitudes for (e) arbitrary numbers and types of external particles. The dual superconformal symmetry is broken at loop level by infrared divergences. However, authors provide evidence that the dual conformal anomaly of (e) the MHV and NMHV superamplitudes is (=) the same and, therefore, their ratio is (=) dual conformally invariant. Authors show this explicitly for (e) the six-particle amplitudes at one loop. Authors conjecture that these properties hold for all, MHV and non-MHV, superamplitudes in N=4 SYM both at (eb) weak and at strong coupling. **Nuclear Physics B Volume 828, Issues 1–2, 21 March 2010, Pages 317–374 Dual superconformal symmetry of scattering amplitudes in N=4 super-Yang–Mills theory J.M. Drummonda DOI: 10.1016/j.nuclphysb.2009.11.022.** By comparison with numerical results in (eb) the maximal Abelian projection of lattice Yang–Mills theory, it is argued that the nonperturbative dynamics of (e) Yang Mills theory can be described by (e) a set of fields that take their values in (eb) the coset space SU(2)/U(1). The Yang–Mills connection is parameterized in (eb) a special way to separate (e&eb) the dependence on the coset field. The coset field is then regarded as (=) a collective variable and a method to obtain its effective

action is developed. It is argued that the physical excitations of (e) the effective action may be (=) knot solitons. A procedure to calculate the mass scale of knot solitons is discussed for (e) lattice gauge theories in the maximal Abelian projection. The approach is extended to (e&eb) the SU (N) Yang–Mills theory. A relation between the large N limit and the monopole dominance is pointed out. **Physics Letters B Volume 458, Issues 2–3, 8 July 1999, Pages 322–330 An effective action for monopoles and knot solitons in Yang–Mills theory Sergei V. Shabanov DOI: 10.1016/S0370-2693(99)00612-7.** It is here the quintessential formulation of consciousness occurs. Like a man standing on the threshold of infinity trying to ponder what lies beyond the veil, which separates the seen from unseen a true seeker tries to lift the individual general ledger to the cosmic general ledger. Like in Darwinian evolution, mind also has to evolve from individual consciousness to cosmic consciousness. In Kaishmiri Shaivism one withdraws from the finite to the infinite, but one also goes on an outward journey from the infinite to the finite, because both the finite and the infinite have an intimate connection. The finite is not seen as unreal, but as a symbol of the infinite. There is no real distinction between them. Those two movements constitute Spanda, a key concept in Kaishmiri Shaivism. Spanda is the pulsation of the Absolute in different phases of being. There are no opposites like subject and object, unity or duality, absolute or relative. They are just different phases of the universal vibration of the Absolute. The goal is to realize or be (=) at once infinite and finite. One does not turn away from appearances (like in Advaita Vedanta), but one realizes that the Absolute manifests all things. Spanda, the eternal pulsation of the Absolute, oscillates between a passion to create and dispassion from the created. Through it the Absolute transform itself into all things and then returns back into the emptiness of its undifferentiated nature. In Kaishmiri Shaivism the Absolute is seen as pure consciousness (=being). The Absolute is an eternal all-pervasive principle, the highest reality, the nature of all entities eternally and blissfully at rest within its own nature. The Absolute is the nature of the Self (and

thus of us all). The Absolute is divine; it is Shiva, the Lord of the Universe. It is full of (e) conscious activity through (e&eb)which it generates the universe, and reabsorbs it into (e&eb) itself at the end of each cycle of creation. Thus we speak of monism, as everything resides within (eb) this one absolute consciousness. It sustains (eb) all things, it embraces (e) all things, and it pervades (e&eb) all things. All things are appearances within (eb) the absolute consciousness, but nevertheless real (in contrast to (e) Advaita Vedanta where (e) appearances are seen as unreal or illusionary). All things appear external (out there, outside ourselves), but they do not have (e) a being on their own. They do not exist as (=) separate entities on their own. Everything is contained within (eb) consciousness. What we see as objects are (=) manifestations of consciousness. The event which constitutes (eb) the universe is always internal events happening within (e&eb) consciousness because (e) their essential nature is consciousness itself. If a physical object were totally material, and independent or external to consciousness, it could never be experienced. The universe and (e&eb) consciousness are two aspects of a whole. The universe is (=) an attribute of consciousness which bears consciousness as (=) its substance. Doctrine of Vibration: S.G.Dyczkowski what you "see" is (=) not always what you "get." Many people mistakenly take their own visions literally without (e) "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in (eb) the physical sense but in a psychological sense. The archetypes of (e) the collective unconscious are arrayed behind (e) the scenes of current worldwide conditions, of (e) crisis and confusion. They mirror (e&eb) our own states back at us, whether we perceive them as (=) such or interpret them plausibly or not (Note the essay has been written in model able form and models are given in various papers.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4\,a_9 \square - b_9 \square 0$$

And

$$b_9^{\,2} \square 4\,a_9 c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9\, in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\, X_9$ ,

$Let\ Y_9 = \ y_9\, X_9$And try the substitution $A_9 = X_9/(2 + A_9\ y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + \ n_9 y_9)$ we have

$$W_9^{\,2}\, y_9 \left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^{\,2}\left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9\, X_9 \left[\frac{X_9 (1 + n_9 y_9)}{M_9^{\,2}(2 + n_9 y_9)^2} + 2\, y_9 \left(n_9 - 1\right) + \frac{X_9^{\,2}(9 + y_9)}{M_9^{\,2}(2 + n_9 y_9)} = 0.\ (A2.12)\right.$$

Since $n_9$ and $A_9$are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9 = (2 + n_9 y_9)$, *the condition*(A2.99?)*can be written as*

$[\frac{X_9(E_9-9)}{E_9^2}]^2[\frac{X_9^2}{M_9^2} - 4\, y_9\ E_9^2 \frac{X_9}{M_9^2} + 4y_9 (n_9 - 2)E_9^3]$ .

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9 / W_9^2)$ can be expressed as

$2\, E_9 [y_9 E_9 + (y_9\ E_9 (9 + y_9)(E_9 - n_9))^{9/2}] > \frac{X_9}{M_9^2} > 2E_9 [y_9 E_9 - y_9 E_9 (9 + y_9)(E_9 - n_9))^{9/2}]$

Concept Of Being From Shaivite Point Of View: In His Famous Work Entitled On The Conservation Of Force (1847), Hermann Von Helmholtz Had Divided The Totality Of The Energy In The Universe Between Potential Energy And Kinetic Energy And Affirmed The Reciprocal Convertibility Of The Two. In 1852, William Thomson Pointed Out That There Exists A Sub-Ensemble Within Kinetic Energy, Heat, Which, Once It Has Been Generated, Is No Longer Entirely Convertible Into Potential Energy-Or Into Any Other Form Of Kinetic Energy. Considering That The (Partial) Reconversion Of Heat Into Labor Is Possible Only In Situations That Present A Disparity In Temperature, And That Heat Tends To Pass From Warmer To Cooler Bodies By Spreading On An Even Temperature Level Through Space, Thomson Concluded That The Universe Tends Towards A Final State Where Any Energetic Transformations, Every Movement And Every Form Of Life Will Cease: The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO: The Ultimate Nature Of Anything Cannot Change. Change Of Any Sort Is Merely Apparent (Vivarta). The World Of Change And Becoming Is A False Super- Imposition (Adhyaropa, Adhyasa) On The Absolute. It Appears So On The Screen Of Individual, Collective Consciousness, Albeit One Cannot Be So Sure Of Cosmic, For There Are Many Reasons That Cosmic Transactions Also Had To Be Subjected To Change

By ParaBrahaman. In Cosmic Terms, The Mistake (Bhranti) Consists Of The Supposition That The Real Brahman Is The Unreal Universe And The Unreal Universe Is The Real Brahman. In Microcosmic Terms, It Is The Mistake Of Falsely Conceiving The Body, Mind Or Even One's Personality To Be The Self. Note We Have Taken Self As Witness Consciousness. Notwithstanding There Is Nothing Wrong In The Conception Of Self As A Dynamical Presupposition, In That Individual Consciousness And Collective Consciousness And Even Cosmic Consciousness Evolutes In To ParaBrahaman. In The Same Way As The Image Of A Snake Is Falsely Superimposed On A Rope, Similarly The Universe Is Falsely Projected Onto The Real Substratum, The Brahman. Ignorance Is Not Merely A Personal Lack Of Knowledge, But A Cosmic Principle. As Such It Is Called "Maya," The Indefinable Factor (Anirvacaniya) That Brings This Mistake In Identity About. The Reality Status Of This Cosmic Illusion Is Also Indefinable: On The One Hand It Is Not Brahman, The Sole Reality; On The Other Hand It Is Not Absolutely Non-Existent Like A Hare's Horn Or The Son Of A Barren Woman. Note Here The Vedantin Avoids Talking About Zero Which We Have Avidly And Assiduously Stressed About. Nothing Is The Factor That Exists. Also See the Remarks about Satre's Being And Nothingness. Brahman Is The Source Of World Appearances Only In The Sense Of Being Their Unconditioned Ground Or Essential Nature. The Universe Is False Not Because It Has No Nature Of Its Own But Because It Does Have One. The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski. Concept Of Being From Shaivite Point Of View: It Is The 'I Am The Body/Mind' Belief That Gives Rise To The 'I Am Not The World' Belief. These Two Beliefs Are (=) Co-Created."Consciousness Projects The Appearance Of The Mind, Body And World By Taking (Eb) The Shape Of Thinking, Sensing And Perceiving."“Attention Is Consciousness With An Object. When The Object Vanishes, Attention Simply Remains What It Always Is Consciousness."“There Is No Purpose To Meditation. The Purpose Is Already Accomplished."“Everything

That Is Experienced Is Experienced By, Through, In And As Consciousness.”“The Seen Cannot Be Separated From Seeing And Seeing Cannot Be Separated From Consciousness.” “**The Reality Of Any Experience Is Not Hidden In The Appearance; It Is Expressed By The Appearance The Transparency Of Things**' By Rupert Spira Book Review By Dennis Waite The Following Is A Review Of Rupert's Book: 'The Transparency Of Things: Contemplating The Nature Of Experience’. There Is An Essay From The Book Here And The Essay On Cézanne May Be Read At The Awakened Eye Website Nature Rejects The Naiveté That Seeks Absolute Truth. We Are Beginning To Realize, Individually And Culturally, That "Realities" Are All Human Constructions. The Task Becomes One Of "Catching Ourselves In The Act" Of Creating Our Own "Reality" From The Flow Of Events. Human Truth Is Always An Engagement Of Mind With Experience. The Challenge Of The Therapist In These Times Of Chaotic Change Is To Validate The Concept That We Don't Have To Fear The Collapse Of What We Think We Are. Strange Attractors: Transference, Holography, And An **Archetype Burke, J. (2003). Strange Attractors: Transference, Holography, And An Archetype (Doctoral Dissertation, Pacifica Graduate Institute, 2003).** When we get to the topic and understanding of the Transcendental Logic, the brilliant and baffling Transcendental Deduction (TD) Recall the two movements just discussed, the one from experience to its conditions and the one from the forms of valid inference to the concepts that we must use in all judging (the Categories). This duality led Kant to his famous question of right (quid juris) (A84=B116): with what right do we apply the Categories, which are not acquired from experience, to the contents of experience? (A85=B117). Kant's problem here is not as arcane as it might seem. It reflects an important question: How is it that the world as we experience it conforms to our logic? In briefest form, Kant thought that the trick to showing how it is possible for the Categories to apply to experience is to show that it is necessary that they apply (A97).TD has two sides, though Kant never treats them separately John Marlado,

accounts, attributes, and ascribes the philosophy of Nishida KitarÅ is lucid, insightful, and deeply informative despite the highly questionable argument that gives the book its structure. In comparison with other, technically more accurate accounts of Nishida's philosophy, it is broader in scope than most but more single-minded in its approach. The analysis is for the most part limited to four books that have been translated into English, out of approximately 16 volumes of philosophical essays in Japanese, but it aims to capture the contours and most important details of Nishida's entire philosophy. To appreciate the achievement of this account, one would need only to read, in English translation or in the Japanese original, any of the volumes with their streams of repetitive, puzzling expressions and circuitous, often dead- end argumentation. Wilkinson's elucidation of details in the passages he samples makes it easier to understand why Nishida is considered Japan's greatest academic philosopher. A critical look at the frame of the book and some of its own claims reveals where, to my mind, it has gone astray. The argument that frames this book is repeated throughout: Nishida's philosophy is an attempt to articulate, in a characteristically Western philosophical manner, the Zen experience he had as a young man. Nishida first made the attempt by using the conceptual frameworks of Western philosophers, and after they failed him he developed his own, largely by recasting Buddhist insights. In the end, a philosophical articulation of the Zen worldview in Western terms or even in Eastern terms recast by conceptual, Western philosophical methods approved impossible. What Nishida attempted is incommensurable with Western philosophy. While the author first offers these statements as hypotheses (p. 2), one premise of his argument that there is such a thing as "the Zen world-view" or "the Zen experience “remains unquestioned; and another that Nishida's pivotal experience was the Zen experienceâis asserted as "beyond question" (p. 151; see also pp. 28, 48, 10, 114).Wilkinson's presentation of Zen in the first chapter is drawn largely from D. T. Suzuki, with help from Suzuki's portage Abe Masao; some comments on time, birth, and death in; and a few other sources. The

result is a caricature that imagines Zen as an invariable, unitary experience. In fact, Dagan's teachings on practice as the manifestation of enlightenment do not sit well with D. T. Suzuki's emphasis on a satori experience, and Suzuki does not enlighten us about what Zen in its various practices has historically been. The author would have benefited from acquaintance with other presentations by Zen teachers, such as that of ShunryÅ Suzuki'sZen Mind Beginner's Mind, with historical accounts of Zen traditions, and with contrasts between DÅgen's view of time and Sam Â· sÄraand that of other Buddhist philosophers. A little research in contemporary Zen scholarship would have shown how implausible is the talk of "nirvana or awareness of mu" (p. 24), "the basic Zen position [that] never varies" (p. 3), "the Zen vision" (pp. 5, 7), "Zen epistemology" , "Zen experience" (p. 59), "the conception of time involved in Zen" (p. 15), and "the fundamental Zen assertion that there is no absolute time" (p. 16). Even if DÅgen's notion that time and being are nondual (p. 16) epitomized most of Zen, Nishida's early notion of "a transcendent, unchanging reality apart from time" (p. 51) would contrast sharply with it. Equally questionable is Wilkinson's contention that Nishida's "pure experience" is equivalent to Zen satori. The notion of pure experience Nishida developed in his first major work may be pivotal for his entire philosophy, but he never claims or implies that any part of his philosophy presupposes an exceptional experience like satori that grounds his convictions and renders them unverifiable to those who have not had the experience, as Wilkinson intimates (pp. 55, 159, and 161). Nishida presents his notions of pure experience and intellectual intuition, the grasp of its unity, and later... Perennialism and Constructivism is studied by Randolph T Dible : experiencer or the observer and the transcendental subjectivity itself Pure objectivity is essentially a deconstruction, apophatic and negativa The encountering of the experiencer or observer—transcendental subjectivity itself—at the foundation of the world leads inevitably to the recognition of pure objectivity as ultimate reality (which can be taken as its ultimate deconstruction, analogous to the apophatic or via

negativa), from which objects derive their value, weight, significance, meaning or objectivity. In this way, pure objectivity can be seen as the supra-self-evident Axiological Axiom, so to speak, even Unconditional Love, in romantic terms. This axiology (value theory) has a structure inverse to the relationship between transcendental subjectivity as the radical unity of pure self-reference and on the other hand, the world of forms, as mere traces (representations, indications) of the unique, original "first distinction" Spencer-Brown speaks of at the foundation of his calculus. That is, all forms (i.e., distinctions, differences) would reduce to being the first distinction, also known as the marked state, which can be called penultimate reality (pure self-reference or transcendental subjectivity: the Spirit which animates us), except that forms are complimentary to their content, which is their objectivity or value, which would reduce to the unmarked state or ultimate reality. It is the incongruity of form (thoughts; Whitehead's "negative prehensions") and value (feelings; Whitehead's "positive prehensions," or my notion of objectivity, meaning and qualia; in short, the non-formal aspects of experience) that holds forms open and keeps them from absolute reduction. This accounts for the brute, concrete persistence of the "functional illusion"-- to use a term from Dzogchen Buddhism-- of the world. Thus this system has an axiology of metaphysical objectivity grounded on the ideal of pure objectivity as the source of all value, meaning and significance, itself the very fecundity of profundity, which is the motive of drawing the distinction in the first place **(Wikipedia, Kant's writings and Logic Of Sense)** Nano rings were created by accident while intending to make quantum dots. They have interesting optical properties associated with excitons and the Aharonov–Bohm effect. Application of these rings used as light capacitors or a buffer includes photonic computing and communications technology. Analysis and measurement of geometric phases in mesoscopic rings is ongoing. Several experiments, including some reported in 2012, show Aharonov-Bohm oscillations in charge density wave (CDW) current versus magnetic flux, of dominant period h/2e through CDW rings up to 85 **μm** in

circumference above 77 K. This behavior is similar to that of the superconducting quantum interference devices (seeSQUID). **Aharonov–Bohm nano rings (Wikipedia) Stability analysis of nano rings would help progressively accentuate the research about the materials division that uses nano rings.** From the profound contemplation and wisdom of the Buddha and Mahabodhisattvas, we know there is no such reality. Instead, egolessness (non-self) is the only path to understand the reality of the deluded life. All existences are subject to the law of causes and conditions. These include the smallest particles, the relationship between the particles, the planets, and the relationship between them, up to and including the whole universe! From the smallest particles to the biggest matter, there exists no absolute independent identity. Egolessness (non-self) implies the void characteristics of all existence. Egolessness (non-self) signifies (eb) the non-existence of permanent identity for self and existence (dharma). Sunyata stresses (eb) the voidness characteristic of (e) self and existence (dharma). Sunyata and Egolessness possess (e) similar attributes. As we have discussed before, we can observe the profound significance of Sunyata from (e) the perspective of inter-dependent relationships. Considering dharma-nature and (e&eb) the condition of nirvana, all existences are immaterial and of (e) a void-nature. Then we see each existence as (=) independent of each other. But then we cannot find any material that does exist independent of (e) everything else. So Egolessness also implies (eb) void-nature! From the observation of all existences, we can infer (eb) the theory of nirvana and the complete cessation of (e) all phenomena. From the viewpoint of phenomena, all existences are so different from (e) each other, that they may contradict (e&eb) each other. They are (=) so chaotic. In reality, their existence is (=) illusionary and arises from (e) conditional causation. They seem to exist on one hand, and yet do not (e) exist on the other. They seem to be united, but yet they are so different to (e) one another. They seem to exist and yet they do cease! Ultimately everything will return to (e) harmony and complete calmness. This is the nature of all existence. It is the

final resting place for all. If we can understand this reality and remove our illusions, we can find this state of harmony and complete calmness. **Teachings in Chinese Buddhism (6) sunyata (emptiness) in the Mahayana context (Wikipedia) Many models have been given concatenating Buddhist doctrines and Vaccum energy. These models essentially throw light on the essence of self(witness consciousness) and the cosmic general ledger (cosmic consciousness) while the imbalance in the gratification deprivation complex always remain unbalanced when it comes to individual general ledger , advocating the evolutionary theory of individual consciousness to cosmic consciousness. It is to be noted that it is not the question of knowing all the transactions but the wisdom of what lies behind the debit- Credit, Debit-Debit, or Credit-Credit principle that is most important for assimilation.** Look at this simple thesis to prove the point stated in the foregoing. Thesis: 'The World Has A Beginning In Time, And In Space It Is Also Enclosed In Boundaries.' Proof: 'For If One Assumes That The World Has No Beginning In Time, Then Up To Every Given Moment In Time An Eternity Is Elapsed, And Hence An Infinite Series Of States Of Things In The World, Each Following Another, Has Passed Away. But Now The Infinity Of A Series Consists Precisely In The Fact That It Can Never Be Completed Through A Successive Synthesis. Therefore An Infinitely Elapsed World-Series Is Impossible, So A Beginning Of The World Is A Necessary Condition Of Its Existence, Which Was The First Point To Be Proved (Kant, Critique Of Pure Reason, Tr. By R. Guyer And A. W. Wood (Cambridge: Cambridge University Press, 1998): B 454, P. 470). As Quoted In The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO"The Cyclical Hypothesis, So Heavily Criticized By Nietzsche (VP II 325 And 334), Arises In This Way." In Fact, Nietzsche Was Not Criticizing The Cyclical Hypothesis But Only The Particular Form Of That Hypothesis Presented In Vogt's Work. All Of Nietzsche's Texts Without Exception Speak Of The Eternal Return As The Repetitio Of The Same Events Within A Cycle Which Repeats Itself

Eternally. If Deleuze's Interpretation Holds That The Eternal Return Is Not A Circle, Then What Is It? A Wheel Moving Centrifugally, Operating A "Creative Selection," "Nietzsche's Secret Is That The Eternal Return Is Selective" Says Deleuze: The Eternal Return Produces Becoming-Active. It Is Sufficient To Relate The Will To Nothingness To The Eternal Return In Order To Realize That Reactive Forces Do Not Return. However Far They Go, However Deep The Becoming-Reactive Of Forces, Reactive Forces Will Not Return. The Small, Petty, Reactive Man Will Not Return. Affirmation Alone Returns, This That Can Be Affirmed Alone Returns, Joy Alone Returns. Everything That Can Be Denied, Everything That Is Negation, Is Expelled Due To The Very Movement Of The Eternal Return. We Were Entitled To Dread That The Combinations Of Nihilism And Reactivity Would Eternally Return Too. The Eternal Return Must Be Compared To A Wheel; Yet, The Movement Of The Wheel Is Endowed With Centrifugal Powers That Drive Away The Entire Negative. Because Being Imposes Itself On Becoming, It Expels From Itself Everything That Contradicts Affirmation, All Forms Of Nihilism And Reactivity: Bad Conscience, Ressentiment..., We Shall Witness Them Only Once. [...] The Eternal Return Is The Repetition, But The Repetition That Selects, The Repetition That Saves. Here Is The Marvelous Secret Of A Selective And Liberating Repetition. There Is No Need To Remind The Reader That Neither The Image Of A Centrifugal Movement Nor The Concept Of A Negativity-Rejecting Repetition Appears Anywhere In Nietzsche's Writings, And Indeed Deleuze Does Not Refer To Any Text In Support Of This Interpretation. Further, One Could Highlight That Nietzsche Never Formulates The Opposition Between Active And Reactive Forces, Which Constitutes The Broader Framework Of Deleuze's Interpretation. For Some Years, Marco Brusotti Has Called Attention To The Fact That Deleuze Introduced A Dualism That Does Not Exist In Nietzsche's Writings. To Be Sure, The German Philosopher Describes A Certain Number Of "Reactive" Phenomena (For Example, In The Second Essay Of The Genealogy Of Morality, § 11, He Talks About

"Reactive Affects" [Reaktive Affekte], "Reactive Feelings" [Reaktive Gefühlen], Reactive Men [Reaktive Menschen]); But These Are Nonetheless The Result Of Complex Ensembles Of Configurations Of Centers Of Forces That Remain In Themselves Active. Neither The Word Nor The Concept Of "Reactive Forces" Ever Appears In Nietzsche's Philosophy. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO Some Parts Are Deleted Due To Spatial Constraints. In the midst of thoughts about the eternal return we find at least two other thematic axes On the one hand, the view of the **world as a constant flux of forces without any goal, law, or rules of becoming. A chaos sive natura de-divinized and de-anthropomorphized which constitutes the "ontological substratum**" of the whole of Nietzsche's reflexions On the other hand, an ensemble of fragments of an anthropologico-sociological character, designing a path of liberation leading to the creation of superior individuals by way of a profound transformation of their instinctual structure. This transformation must be achieved by a practice of solitude and internal struggle towards the liberation from the ancient representations of the world and from the incorporated herd values. For an analysis of these thematic perspectives, see Paolo D'Iorio, La linea e il circolo. Cosmologia e filosofia dell'eterno ritorno in Nietzsche (Genova: Pantograf, 1995): 233-322. With most of the parameters quantified, with assumptions there must be picturesque view of the self and non self, space and time, science and philosophy with the concatenational consummations. After The Chapter "On Redemption," Where Zarathustra Dares Not Expose His Doctrine, The Eternal Return Begins To Be Enunciated In Part Three Of The Work. In The First Place, It Is The Dwarf Who Formulates It In The Chapter "On The Vision And The Riddle." Facing The "Gate Of The Instant" Which Symbolizes The Two Infinities That Stretch Towards The Past And The Future, The Dwarf Whispers: "All Truth Is Crooked, Time Itself Is A Circle" The Dwarf Represents The Spirit Of Gravity, And He Embodies The Herd Morality, "The Belittling Virtue" Which Is The Title Of Another Chapter From Part III. The Dwarf Can

Endure The Eternal Return Without Great Difficulties Because He Has No Aspirations; Unlike Zarathustra He Does Not Wish To Climb The Mountains That Symbolize Elevation And Solitude. In Two Unpublished Notes, From The Summer And The Fall Of 1883, Nietzsche Writes: The Doctrine Is At First Favored By The Rabble, Before It Gets To The Superior Men. The Doctrine Of Recurrence Will First Smile To The Rabble, Which Is Cold And Without Any Strong Internal Need. It Is The Most Ordinary Of Life Instincts, Which Gives Its Agreement First. Hence, The Content Of The Doctrine Is The Same, But Whereas The Dwarf Can Endure It (Because He Interprets It According To The Pessimistic Tradition For Which "Nothing Is New Under The Sun"), Zarathustra, Who Is The "Advocate Of Life" Regards The Eternal Return As The Strongest Objection To Existence, And As The Rest Of The Dream Suggests, He Does Not Yet Succeed In Accepting It After The Vision At The Gate Of The Instant, The Chapter Is Brought To An End By The Enigma Of The Shepherd. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO .How does the memory affect the GD complex or for that matter the gratification deprivation of an individual or a collective identity? Non-Markovian local-in-time master equations give a relatively simple way to describe the dynamics of open quantum systems with memory effects. Despite their simple form, there are still many misunderstandings related to the physical applicability and interpretation of these equations. Here, we clarify these issues both in the cases of quantum and classical master equations. Further introduction of the concept of a classical non-Markov chain signified through negative jump rates in the chain configuration. **E-M Laine et al 2012 J. Phys. B: At. Mol. Opt. Phys. 45 154004 doi:10.1088/0953-4075/45/15/154004 Local-in-time master equations with memory effects: applicability and interpretation.** Supposing That There Were Indeed An "Energy Of Contraction" Constant In All Centers Of Force Of The Universe, It Remains To Be Explained Where Any Difference Would Ever Originate. It Would Be Necessary For The Whole To Dissolve Into An Infinite Number Of

Perfectly Identical Existential Rings And Spheres, And We Would Therefore Behold Innumerable And Perfectly Identical Worlds COEXISTING [Nietzsche Underlines This Word Twice] Alongside Each Other. Is It Necessary For Me To Admit This? Is It Necessary To Posit An Eternal Coexistence On Top Of The Eternal Succession Of Identical Worlds? The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. Science is usually portrayed as dehumanizing. Brave New World epitomizes this fear. "The more we understand the world, the more it seems completely pointless" (Steven Weinberg). Certainly science can seem chilling when conceived in the abstract as a metaphysical world-picture. We may seem to find ourselves living in a universe with all the human meaning stripped out: Participants in a soulless dance of molecules, or harmonics of pointlessly waggling superstrings and their braneworld cousins. Nature seems loveless and indifferent to our lives. What rights have we to be happy? **(Brave New World: Aldous Huxley).** Religion can afford to be dogmatic but science cannot. Look at this passage which counters both scientists and philosophers alike: However, Hartmann Objects That The Regressive Movement Postulated By Schopenhauer Is Possible Only In Thought: It Remains Nothing More Than An "Ideal Postulate" With No Real Object And Which "Does Not Teach Us Anything About The Real Process Of The World That Unfolds In A Movement Contrary To This Backwards Movement Of Thought" (Hartmann, Philosophie Des Unbewussten, Third Edition (1871): 772). Hartmann Affirms That If Unlike Schopenhauer One Admits The Reality Of Time And Of The World Process, One Must Also Admit That The Process Must Be Limited In The Past And Therefore That There Must Be An Absolute Beginning. In Hartmann's Mind, Failure To Do So Would Result In Positing The Contradictory Concept Of An Accomplished Infinity: "The Infinity That From The Point Of View Of Regressive Thinking, Remains An Ideal Postulate, Which No Reality May Correspond To, Must, For The World, Whose Process Is, On The Contrary, A Progressive Movement, Open Up To A Determinate Result; And Here The Contradiction Comes To Light" (Hartmann

(1871): 772). What Really "Comes To Light" In This Passage Is The Fact That Hartmann Does Not Provide A Demonstration But A Petitio Principii. Indeed, The Concept Of The World Process Analytically Contains The Concept Of A Beginning Of The World. In All Rigors, It Is Therefore Impossible To Demonstrate These Concepts With Reference To Each Other. Secondly, Hartmann's View That One Is Bound To Accept The Reality Of The World Process Even If One Rejects The Ideality Of Schopenhauer's Time Is Mistaken. Hartmann Believes That If Time Is Real There Must Be A World Process With Both An Absolute Beginning And An Absolute End. Without Any Justification, Hartmann Jumps From Schopenhauer's Negated Time To Oriented Time. The Eternal Return: Genesis And Interpretation BY PAUL D'IORIO. In the last years several theoretical papers discussed if time can be an emergent property deriving from ) quantum correlations. Here, to provide an insight into how this phenomenon can occur, authors present an experiment that illustrates Page and Wootters' mechanism of "static" time, and Gambini et al. subsequent refinements. A static, entangled state between a clock system and the rest of the universe is perceived as evolving by internal observers that test the correlations between the two subsystems. They implement this mechanism using an entangled state of the polarization of two photons, one of which is used as a clock to gauge the evolution of the second: an "internal" observer that becomes correlated with the clock photon sees the other system evolve, while an "external" observer that only observes global properties of the two photons can prove it is static.Phys. Rev. A 89, 052122 (2014)DOI:10.1103/PhysRevA.89.052122 arXiv: 1310.4691 [quant-ph] **Time from quantum entanglement: an experimental illustration Ekaterina Moreva.** Time is an entanglement .When the new ideas of quantum mechanics spread through science like wildfire in the first half of the 20th century, one of the first things physicists did was to apply them to gravity and general relativity. The results were not pretty. It immediately became clear that these two foundations of modern physics were entirely

incompatible. When physicists attempted to meld the approaches, the resulting equations were bedeviled with infinities making it impossible to make sense of the results. Then in the mid-1960s, there was a breakthrough. The physicists John Wheeler and Bryce DeWitt successfully combined the previously incompatible ideas in a key result that has since become known as the Wheeler-DeWitt equation. This is important because it avoids the troublesome infinites—a huge advance. But it didn't take physicists long to realise that while the Wheeler-DeWitt equation solved one significant problem, it introduced another. The new problem was that time played no role in this equation. In effect, it says that nothing ever happens in the universe, a prediction that is clearly at odds with the observational evidence. This conundrum, which physicists call 'the problem of time', has proved to be thorn in flesh of modern physicists, who have tried to ignore it but with little success. Then in 1983, the theorists Don Page and William Wootters came up with a novel solution based on the quantum phenomenon of entanglement. This is the exotic property in which two quantum particles share the same existence, even though they are physically separated. Entanglement is a deep and powerful link and Page and Wootters showed how it can be used (e) to measure time. Their idea was that the way pair of entangled particles evolves is a kind of clock that can be used to measure change. But the results depend on how the observation is made. One way to do this is to compare the change in the entangled particles with an external clock that is entirely independent of the universe. This is equivalent to god-like observer outside the universe measuring the evolution of the particles using an external clock. In this case, Page and Wootters showed that the particles would appear entirely unchanging—that time would not exist in this scenario. But there is another way to do it that gives a different result. This is for an observer inside the universe to compare the evolution of the particles with the rest of the universe. In this case, the internal observer would see a change and this difference in the evolution of entangled particles compared with everything else is an important a measure of time. This is

an elegant and powerful idea. It suggests that time is an emergent phenomenon that comes about because of the nature of entanglement. And it exists only for observers inside the universe. **Time is an entanglement** (Wikipedia). Danish religious philosopher Soren Kierkegaard carried out a systematic critique of the pretensions of reason and an abstract rationalism which he believed that the modern age was nurturing. Condemning reflection as a "danger" that ensnares people in logical delays and machinations, Kierkegaard compared it to a prison. Reflection is for him a form of captivity, a bondage which "can only be broken by [passionate] religious inwardness" (1978: 81). **Reflection seduces individuals into thinking its possibilities are "much more magnificent than a paltry decision"** (1978: 82). It leads them to act **"on principle,"** to dwell on the deliberation of the context of their actions and the calculation of their worth or outcome. Kierkegaard argues that this drives away feeling, inspiration, and spontaneity, all of which are crucial for true inner being and a vital relation to God. For Kierkegaard, as Nietzsche would later agree, genuine inner being (and culture) is characterized by the tautness and tension of the soul which characterizes passionate existence. But the "coiled springs of life relationships ... lose their resilience" in reflection (1978: 78) and "everything becomes meaningless externality, devoid of [internal] character" (1978: 62). Kierkegaard thus contributes to the development of an irrationalist tradition that has echoes in some later postmodern thought. Kierkegaard might well have agreed with his contemporary Fyodor Dostoyevsky, who wrote: "An intelligent [reflective] man cannot seriously become anything ... excessive consciousness is a disease" (1974: 3, 5). In an age overtaken by rules and regulations, genuine action -- which Kierkegaard assumes to be subjective and spontaneous -- is frustrated at every turn. Complaining that we are too "sober and serious" (1978: 71) even at banquets, Kierkegaard bemoans the fact that even suicides are premeditated (1978: 68)! "That a person stands or falls on his actions is becoming obsolete; instead, everybody sits around and does a brilliant job of bungling through with the aid of some reflection

and also by declaring that they all know very well what has to be done" (1978: 73). Thus, it is passion, not reflection, that guarantees "a decent modesty between man and man [and] prevents crude aggressiveness" (1978: 62). "Take away the passion and the propriety also disappears" (1978: 64). The ambiguity in the word "passion" may cause some confusion here. To say that the age and its individuals are "passionless" is not to say there are no emotions whatsoever, but rather that there is no true spiritual inwardness and depth, no intensively motivated action and commitment. It suggests that passion exists only in a simulated, pseudo-form, "the rebirth of passion" through "talkativeness" (1978: 64). "Chattering" for Kierkegaard gets in the way of "essential speaking" and merely "reflects" inconsequential events (1978: 89-99). Hence, in "the present age," emotions -- which in fact are all too pronounced -- have been transformed into negative forces. Anticipating Nietzsche's genealogy of the "slave revolt" in morality, Kierkegaard claims that the "enthusiasm" of the prior age of Revolution, a "positively unifying principle," has become a vicious "envy," a "negatively unifying principle" (1978: 81), a leveling force in its own right insofar as those lacking in talent and resources want to tear down those who have them. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS By Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html.** In November 2011, the United States government argued before the US Supreme Court that it wants to continue utilizing GPS tracking of individuals without first seeking a warrant. In response, Justice Stephen Breyer questioned what this means for a democratic society by referencing Nineteen Eighty-Four. Justice Breyer asked, "If you win this case, then there is nothing to prevent the police or the government from monitoring 24 hours a day the public movement of every citizen of the United States. So if you win, you suddenly produce what sounds like 1984... "[65]. What is to be said in unmistakable terms is that when it comes to national security and innocent citizens there is question of compromise even it be an infringement and encroachment of so called

privacy which has been made a word to subterfuge the activities that are far said to be fair. Quantum theory has provoked intense discussions about its interpretation since its pioneer days. One of the few scientists who have been continuously engaged in this development from both physical and philosophical perspectives is Carl Friedrich von Weizsäcker. The questions he posed were and are inspiring for many, including the authors of this contribution. Weizsäcker developed Bohr's view of quantum theory as a theory of knowledge

## ABSOLUTE STABILITY

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_9$ ) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9 - 9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2 - 2y_9\frac{(E_9 + n_9^2 y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right) - 4y_9(E_9 - n_9)$$

And has the solution

$$\frac{X_9}{M_9^2} > \frac{E_9}{(E_9 - 9)}[y_9(E_9 + n_9^2 y_9) + [y_9^2 (E_9 + n_9^2 y_9)^2 + 4y_9(E_9 - 9)(E_9 - n_9)]^{9/2}]$$

Given $E_9 > n_9$ this lower limit to $X_9 / M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9 / M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$ in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9 < n_9$ $(ie., 2A_9 M_9^2 < X_9 - Y_9)$ then (A2.9)and (A2.8)show that $Q_9$ must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2} > 2E_9[y_9 E_9 - (y_9 E_9 (9 + y_9)(E_9 - n_9)]^{9/2}]$$

And

$$E_9 > n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2 /2M_9 (X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$ and $X_9$ and $Y_9$ by $W_9^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9$ *and* separate (A3.1)into real and imaginary parts

$$Z_9^2 + (M_9 + V_9) Z_9 \sin(Z_9 T_9) - (U_9 M_9 + Y_9) + (M_9(U_9 + V_9) - X_9)$$

$$\mathrm{X}\cos(Z_9 T_9) - 2M_9 V_9 \cos^2(Z_9 T_9) + M_9 V_9 = 0$$

$$(M_9 + U_9) Z_9 - (M_9 + V_9) Z_9 \cos(Z_9 T_9) + (M_9(U_9 + V_9) - X)$$

$$\mathrm{X}\sin((Z_9 T_9) - M_9 V_9 \sin(Z_9 T_9) = 0$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 = 1 - \cos Z_9 T_9$and $A_9 = U_9 - V_9$we obtain

$$4 M_9^2 V_9^2 W_9^4 + 2 M_9 V_9 W_9^3[(M_9 - V_9)^2 + A_9(M_9 + V_9) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2[M_9 - V_9)^2(2 A_9 M_9 + Y_9 - X_9) + (2A_9 M_9 - X_9)^2 - A_9(M_9 - V_9)(2A_9 M_9 - X_9)$$

$$+ M_9 V_9 (A_9^2 + 4(X_9 + Y_9))]$$

$$+ W_9[A_9^2 + 2X_9)(2 A_9 M_9 - X_9) + 2A_9 Y_9(M_9 + V_9) - M_9 A_9^3] + A_9^2(X_9 M_9 - Y_9) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9\left(2, U_9 V_9\right) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting cos $(Z_9 T_9) = -1\ in\ (A3.3)$

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**LEN NUNNEYS**

**Absolute Stability in Predator- Prey Models**

**APPLIED TO ACCETUATION ATTRITION MODELS**

**MDDULE NUMBERED SEVEN**

An equilibrium of a time- lagged model is said to be absolutely stable if it remains locally stable regardless of the length of the time delay, and it is argued that the criteria for absolute stability provide a valuable guide to the behavior of models. For example, it is sometimes assumed that time delays have a limited impact until they exceed the natural time scale of a system; here it is stressed that under some conditions very short time delays can have a marked (and often maximal) destabilizing effect. Consequently it is important that our understanding of population dynamics is robust to the inclusion of the short time dealys present in all biological systems.The absolute stability

criteria for absolute stability rather than using criteria which depend upon the details of a time delay is that time delays are unlikely to be constant. For example, a time delay due to debitability inevitably varies between systems and the mean may itselfvary over time. Here it is shown that the criteria for absolute stability are generally robust in the presence of distributed delaysand of varying delays. The analysis presented is based upon a general debit credit system model and it is shown that absolute stability can be expected under a broad range of parameter values whenever the time delay is due to the maturation time of either the debit or the credit of both. This stability occurs because of the interaction between delayed and undelayed dynamic features of the model. A time-delayed process, when viewed across all possible delays, always reduces stability and this effect occurs regardless of whether the process would act to stabilize or destabilize an undelayed system. Opposing the destabilization due to a time delay and making absolute stability a possibility are a number of processes which act without delay. Some of these processes can be identified as stabilizing from the analysis of undelayed models (for example, the non performance assets in a box numerical response of credit types).

**Morphological And Epistemological And Syllogistic Models For Which Nunny's Model Could Be Applied And Has Been Applied:**

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production**

**Charge Asymmetry, $\Delta Y \equiv Yt - Yt^{-}$, Reflects The Asymmetry In $Tt^{-}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles

ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic

quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation

between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to

nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of

Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first

impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of

another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease. Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation

of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces

always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the

thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of

thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity

and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI

Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs” Outcomes—

And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities

Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

$$\frac{dN}{dt} = NF\ (R) - \ NM(R)$$

(1)

$$\frac{dR}{dt} = B(R) - \ D(R) - \ NG(R)$$

Where N and R define the credit type and debit type (resource) levels. The percapita birth and death rates of the credit types confederation are functions of debit types and are defined by the increasing function F and the decreasing function M, respectively. The debit credit rates of the debit confederation are defined at the population level by the functions B and D. D is an increasing (unregulated birth) or decreasing (regulated birth) depending upon where the prey equilibrium lies in relation to the peak of the recruitment curve. Finally the capture rate or functional response of each credit confederation is defined by G(R).

Using the undelayed model (1), it can easily be shown that an equilibrium will be locally stable if and only if

$$\frac{dD}{dR} + N\frac{dG}{dR} - \frac{dB}{dR}$$

$> 0.$ (2)

This condition, like all other stability conditions given in this paper, is evaluated at the equilibrium. The potential importance of debit type self- regulation ($dB/dR < 0$) and of the accelerating part of the type 3 functional response ($dG/dR < G/R$) in promoting the stability of debit credit models generalized by (1) is well known (for a review see Murdoch and Oaten, 1995); however, in time- lagged models an important difference between these two factors must be borne in mind: debit types self- regulation may act as a delayed term but the effect of predation is always immediate.

The necessary and sufficient condition for local stability, defined by (2), depends only upon components of the debit consideration dynamics. The presence of the credit type consummation is felt only indirectly via its regulation of debit distraction ($N\ dG/dR$). It will be seen that the left-hand side of (2), which will be reffered to as the total credit type regulation, continues to play a central role in determining local stability in the presence of time delays, but a delay in the credit confederation has the effect of incorporating features of the credit dynamics into the stability criteria. First, however, consider the simplest case, the effect of a single delay due to debit maturation time, as in cased us bills under both LC and NON LC.

## DELAY DUE TO DEBIT CONFEDERATION APPLICABILITY

Assuming that the time taken for the prey (or resource) to mature is T, and that this is long compared to the maturation period of the predator; then, model (1) can be rewritten as

$$\frac{dN(t)}{dt} = N(t)\big[F\big(R(t)\big) - M\big(R(t)\big)\big]$$

(3)

$$\frac{dR(t)}{dt} = B\big(R(t-T)\big) - D\big(R(t)\big) - N(t)G\big(R(t)\big).$$

It is assumed that debit raised at time $(t-T)$ determine at time $t$ and enter the next confederation dynamics. Furthermore because of differences in ecology, size, or other factors, it is assumed that the are not subject to significant and do not compete with the adult prey for food.

It has been shown (Nunney, 9985a) that the model (3) will retain local stability regardless of the length of the time delay, T, provided that (at the equilibrium)

$$\frac{dD}{dR} + N\frac{dG}{dR} > \left|\frac{dB}{dR}\right|. \qquad (4)$$

Hence the inequality (4) defines the necessary and sufficient condition for absolute stability. This condition can be viewed as two distinct criteria:

We can write (3) as:

$$\frac{d(N(t)_9}{dt} = [N(t)]_9[(F_9)(R(t)]_9 - [M]_9[R(t)]_9$$

$$\frac{d[R(t)]_9}{dt} = [B]_9\,[(R)_9(t-T_9)] - [D]_9[(R(t)]_9 - [N(t)]_9\left[G\big(R(t)\big)_9\right]$$

Each term can be written in the following equations:-

$$\frac{[N(t)]_9}{dt} = a_9[N(t)]_9 - a_9[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_9 = a_3[N(t)]_3 - a_4[N(t)]_9$$

$$\frac{d}{dt}[N(t)]_3 = a_3[N(t)]_9 - a_6[N(t)]_3$$

Each and above term can be written in the above manner and haimovici model can be applied.

First, the total debits regulation $(dD/dR + NdG/dR - dB/dR)$must be positive, as is required in the undelayed system $[see(2)]$. This condition constrains $dB/dR$ frim being too positive.Second, a new feature, the effective debit regulation $(dD/dR + NdG/dR + dB/dR)$, must also be positive. The effective debit regulation treats the delayed effect as potentially destabilizing. Strong regulation of recruitment, when subject to certain time delays, can produce an entirely inappropriate nonregulatory response. Hence this condition constrains $dB/dR$ from being too negative. It also illustrates an important rule: when viewed over a range of potential delays, it is the magnitude rather than the sign of a delayed process which determines stability.

In an undelayed system, a large negative value of $dB/dR$ reflects strong self-regulation and has a powerful stabilizing effect. Self-regulation at the lowest level is often employed to stabilize model systems. The condition (4) suggests that the stabilizing role of self-regulation via recruitment may be limited. This limitation would be trivial if only very long delays minimized the role of self –regulation (since few organisms could have such an exceedingly long maturation as in debit however, the beneficial effect of debit self –regulation is lost very rapidly and time lags much less than the natural period of the debit –credit system can cause a significant loss of stability. This effect can be seen in fig.9b, which represents a system with self-regulation acting through recruitment ($dD/dR$<0). The regulation is strong enough to put the system at the limit of absolute stability. By contrast, the example shown in fig. 9a is at the

**TIME DELAY**

FIG. 9. The relationship between stability and the length of a debit maturation delay. Stability is measured by changing the sign of the real part of the dominant eigenvalue. The time delay axis is scaled to the natural period, which is defined as $2\pi/b$ when the undelayed system has a domunent eigenvalue of $-a \pm ib$. The examples illustrated are both at the extreme limit of absolute stability. (a) The total debit regulation equal to nzero, such as in the Lotka-Volterra predator-prey model. Stability is lowest when there is no delay. (b) The effective debit regulation equal to zero. Stability is lowest with a time delay of roughly half the natural period. Parameters: In each case $M = 2$, $(NGdF/dR - NGdM/dR)2 = dD/dR + NdG/dr = 0.9$, but $dB/dR$ is varied: (a) 0.1, (b)-0.1.

Limitations of Predator prey analysis in immense. When we realize that world itself is Debit-Credit world and sometime credit-Credit world (See an earlier paper entitled the same nomeneclature.

Opposite limit of absolute stability. The result is a specific The Doctrine of Vibration awareness of an object of the form: 'here this is such.' While individual appearances do not lose their separate identity {svarupabheda) when they rest on a common basis, even so the particular object which appears according to its own characteristics (svalak$ana) is an individual reality in its own right. It is a different kind of appearance characterised by its association with the appearance of the specific location and time in which it is made manifest. The form of our experience is thus 'I now see this here'. But when we perceive each particular constituent appearance separately, each assumes a separate fixed function. Abhinava cites the following colourful example to illustrate how the various combinations of appearances account for the variety of experience: Thus even though the appearance of the beloved may manifest externally, it is as if far away in the absence of another appearance, namely, that of 'embracing'. So when the [appearing of the beloved] is associated with another appearance [namely that of 'far away'] the power (arthakriya) it formerly had of giving pleasure appears as its contrary. The form our experience assumes depends, not only on the nature of the object perceived, but also on personal factors entirely peculiar to ourselves. This theory explains this in two ways. In one sense, the object remains the same, but one or other of its constituent appearances

comes to the fore according to the inclinations of the perceiver. From another point of view, we can say that the perceived object is different for each perceiver according to the difference in the prominent appearance manifest to him. Abhinava, citing as an example a golden jar, illustrates how the same object appears differently to different perceivers according to the use they wish to make of it and to their state of mind: When a person who is depressed and feels that there is nothing [of value for him in the world] sees the jar, he merely perceives the appearance 'exists' [in the form of the awareness that] 'it is\ He is not conscious of any other [of its constituent appearances] at all. An individual who desires to fetch water [perceives] the appearance 'jar'. The man who simply wants something that can be taken somewhere and then brought back [perceives] the appearance 'thing'. The man who desires money [perceives] the appearance 'gold'. The man who desires a pleasing object [perceives] the appearance 'brightness' while he who wants something solid sees the appearance 'hardness'. These 'atomic events' or appearances emerge from the pure subject's consciousness and combine together to form a total event at each moment. Integral Monism of Kashmiri Saivism Daily life (vyavahara) goes on by virtue of this ever renewed flux of appearances. They are connected together and work towards a single unified experience because they appear within the field of consciousness of the universal subject. The aggregate of appearances arises in the [supreme] subject as do [sprouts in] a rice field. Even though each sprout germinates from its own seed, they are perceived as a collective whole. Appearances rest in this way within the universal subject. 'External- ity' is itself another appearance; it arises from a distinction between appearances and the individual subject. So, although all manifestation always occurs within the subject, it appears to be external due to the power of Maya which separates the individual subject from his object. This split must occur for daily life to be possible. Only externally manifest appearances can perform their functions; when they are merged within the subject and at one with him, they cannot do so. Daily life proceeds on the basis of the operation and withdrawal of the conditions necessary for fruitful action to be possible. Appearance in this sense represents the actualisation of a potential hidden in consciousness made possible by virtue of its dynamic, Spanda

nature which is both the flow from inner to outer and back as well as the power that impels it. The emergence from, and submergence into, pure consciousness of each individual appearance is a particular pulsation (visesaspanda) of differentiated awareness. Together these individual pulsations constitute the universal pulse (samanyaspandd) of cosmic creation and destruction. Thus, every single thing in this way forms a part of the radiant vibration {sphuratta, sphurana) of the light of absolute consciousness. Light and Awareness: The Two Aspects of Consciousness Absolute consciousness understood as the unchanging ontological ground of all appearing is termed 'Prakasa'. As the creative awareness of its own Being, the absolute is called ' Vimarsa '. Prakasa and Vimarsa — the Divine Light of consciousness and the reflective awareness this Light has of its own nature — together constitute the all-embracing fullness (purnata) of consciousness. The Recognition (pratyabhijfia) school of Kashmiri Saivism develops this concept of the absolute which finds its fullest expression in Utpaladeva's Stanzas on the Recognition of God. Even though neither of these two key terms appear in the Stanzas on Vibration or the Aphorisms of Siva, they recur frequently in their commentaries. Thus, although the original formulation of the Doctrine of Vibration differs from the theology of Recognition in this respect, it was extended in the course of its development to accommodate this concept of the absolute as well. This was possible, and quite justified, insofar as the absolute understood in Pratyabhijfia terms does not, as we shall see, differ essentially from that of the Spanda school. We can, as Kashmiri Saivites themselves have done, explain one in terms of the other. The Doctrine of Vibration Prakasa: The Light of Consciousness Prakasa is the pure luminosity' (bhdna) or 'self-showing' that constitutes the essence and ultimate identity (atman) of phenomena. That things appear at all is due to the light of consciousness, and their appearing (avabhasana) is itself this Light which bestows on all things their evident, manifest nature. Established in the light of consciousness everything appears there according to its own specific nature (svabhava). Anything that supposedly does not rest in this Light is as unreal as a sky-flower. 3 Thus, according to Rajanaka Rama, unlike the light of the sun, or any other light, this Light not only makes all things apparent, it is also their ultimate source. 4 Full of its

divine vibration the Light makes all things manifest and withdraws them into itself. This supra- temporal activity characterises it most specifically; devoid of it, it would be no better than an inert physical phenomenon. At the same time, this light is the conjunction (slesa) or oneness (aikdtmya) of its countless manifest forms, 6 and the collective whole (Sampina'ana) of all the categories of existence. The universe is nothing but the shining of the Light within itself. It is the radiant vibration (sphuratta) of this Light, the state (avast hand) in which consciousness becomes manifest. Although the Light shines as all things at all times and hence also makes their diversity manifest, 9 penetrating each object individually as well as collectively, it is not totally 'merged' (magna) or identified with the object so as to suffer any division within itself. Our experience of any object is of the form: T see this': it is not itself an object, but the manifest form the object assumes as a luminous principle of experience. The Light is ever revealed and can never be obscured; objectivity can never cast a shadow on the light of consciousness. The Stanzas on Vibration declare: That in which all this creation is established and from whence it arises is nowhere obstructed because it is unconditioned by [its very] nature. This Light is the highest reality (Paramartha). lity with $dB/dR$>0. The regulation of debit recruitment is insuffient to stabilize the undelayed system which, like the Lotka-Volterra model, is neutrally stable. However, given a short time delay in recruitment, the initially neutral equilibrium becomes strongly stabilized. The same effect can also stabilize an initially unstable system (Hastings, 1983). This somewhat unexpected role of a time delay increasing stability is discussed in detail elsewhere (Nunney, 1985b).

The criterion of absolute stability is not appropriate for examining the details of the interaction between a time delay and stability. Its value is that it provides a sufficient condition for stability which is extremely robust.

**DELAY DUE TO CONFEDERATION DEBIT TYPE MATURATION CREDIT TYPE CONSOLIDATION MATURATION**

If a predator takes a time $T$ to mature, and if $T$ is long compared to the maturation time of the prey, then the model (1) can be rewritten as

$$\frac{d[NCt]_9}{dt} = N_9(t-T_9)F_9(R_9(t-T_9) - N_9(t)M_9((R_9(t) \quad (5)$$

$$\frac{dR_9(t)}{dt} = B_9\big(R_9(t) - D_9(R_9(t)\big) - N_9(t)G_9(R_9(t))$$

This form of delay relies upon the implicit assumption the survival of is not dependent upon the credit level and the these initial credits have no significant impact upon the level of the debit. As in the previous example of a prey delay, this assumption serves to highlight the difference between delayed and undelayed models by preventing the initial debits from having any dynamic role.

Nunney (1985a) has shown that, as $T$ approaches infinity, the system (5) will be stable if

$$\frac{dD_9}{dR_9} + \frac{N_9 dG_9}{dR_9} - \frac{dB_9}{dR_9} > \frac{N_9 G_9}{2M_9}\left[\frac{dF_9}{dR_9} + \frac{dM_9}{dR_9}\right], \quad (6)$$

(Evaluated at the equilibrium) provided that $dM_9/dR_9$ is not too large. We can usefully simplify this form. The left-hand side of (6), the total debit regulation, can be normalized by $M_9$ This term will be defined by $P$.Further, we define two terms, also noemalized by $M$, which encapsulate effects acting on the credit condition, one derived from the time-delayed or "legged" processes ($L_9$) and the other derived from the non-delayed or "immediate" processes($I_9$):as in payment of usane bills bin according with terms and conditions of the letter of credit.

The appearance may be an illusion, but the illusion itself is real. There is an illusion. It has Reality." He refers to the rope-snake metaphor and says that: "We do not see anything new. We see in a new way." "So that we know that nature is real, that there is something present, that there is a reality to it, even if everything that appears to us is insubstantial and fleeting." **The Transparency of Things' by Rupert Spira Book Review by Dennis Waite The following is a review of Rupert's book: 'The Transparency of Things: Contemplating the Nature of Experience'. There is an essay from the book here and the essay on Cézanne may be read at the awakened**

**eye website For models see Mitigation Of Beam-Induced Backgrounds, Multi-Particle Azimuthal Correlations, Search For Extra Dimensions, Multi-Higgs-Boson Cascade, Top Quark Pair (Tt) Production Charge Asymmetry, $\Delta Y \equiv Yt - Yt^{-}$, Reflects The Asymmetry In $Tt^{-}$ Production, Ultra-Relativistic Electrons And Other Letters Of Interest: A Synecdochal Syncretism: Et Lux In Tenebris Lucet: Light Shines In The Darkness Models.** What then is the relation of Higgs boson and consciousness? Functionalist approaches generally assume that conscious experience appears as a novel property at a critical level of computational complexity. On the surface this would seem to deal with issues 1 and 3, however a conscious threshold has neither been identified nor predicted, and there are no apparent differences in electrophysiological activities between non-conscious and conscious activity. Regarding the nature of experience (why we are not unfeeling "zombies") functionalism offers no testable predictions. Problem 2) of 'binding' in vision and self is often attributed by functionalists to temporal correlation (e.g. coherent 40 Hz), but it is unclear why temporal correlation per se should bind experience without an explanation of experience. As functionalism is based on deterministic computation, it is also unable to account for Penrose's proposed non-computability (4), or free will (5). Something may be missing. To address these issues, various proposals have been suggested in which macroscopic quantum phenomena are connected to the brain's known neural activity. For the problem of unitary binding, Marshall (1989) suggested that coherent quantum states known as Bose-Einstein condensation occurred among neural proteins (c.f. Penrose, 1987; Bohm and Hiley, 1993; Jibu and Yasue, 1995). Pre-conscious to conscious transitions was identified by Stapp (1992) with collapse of a quantum wave function in pre-synaptic axon terminals (c.f. Beck and Eccles, 1992). In another proposal, protein assemblies called microtubules within the brain's neurons are viewed as self-organizing quantum computers ("orchestrated objective reduction - Orch OR" e.g**. Penrose and Hameroff, 1995; Hameroff and Penrose 1996a; 1996b; c.f. Hameroff 1997; 1998a; 1998b; 1998c; 1998d). Stability analysis of such theories would certainly provide a gateway for further bringing in the "reality" in space time.** String theory suggests that the graviton

particle would not be detected from a particle collision because it would have "jumped" into another dimension. So, if the graviton exists and the theory is correct, there should be a "gap" in the energy profile of all the particles ejected from the collision. To understand this, think of breaking the pack of balls in a game of pool. The total energy of all the balls after the collision will depend on the force with which the cue ball was hit. So, after the collision, the energy of all the balls is measured and added-up to make sure it is the same as the energy of the cue ball. BUT - imagine that there is some energy missing, because one of the pool balls has "gone missing" - into a higher dimension. (POSTED BY THEREALJEFFHAL**) Talbot says** "If the universe is a hologram, in some sense it suggests that there may be two very drastically different levels of reality: the concrete reality that we see when we look at [things]... and at some deep level there's a level of reality where everything dissolves into an ocean of energy that is holographically interconnected.". A theory of everything must encompass all the variations however contradictory it may be. Nor need any theory be sent in to grand limbo of oblivion and hibernation or a differentiated groundless less ness because there are some irregularities which are like audit discrepancies need to be rectified. Although the language of religion and science seem different, both are talking about the central order of nature relative to DNA and information, as is Peirce, Petrus, Bacon, and pharaonic Egypt. This path-ordering, active information suggests that the central pattern or law of our holographic cosmos may be modeled on the biophysical battleground of gene regulation in bacteriophage Lambda, one of nature's most efficient and highly evolved mechanisms identified by experimental praxis. Egyptian texts support that the Lambda lifestyles of lysogeny and lysis can be understood as two paths to two living systems, photosynthesis and its reversal chemiluminescence, the cosmic key for evolution of mind and continuity. The texts support that the Lambda genome is the world-heart of two ways. The viral lifestyle of lysogeny controlled by cI protein results in matter, classical spacetime, our world of projected shadow, photosynthesis, sexual genesis or vertical gene transfer, mind in the human body, and the imperfect continuum. In contrast, the viral lifestyle of lysis controlled by cro protein and lactose metabolism results

in transformation to energy, the quantum, chemiluminescence, asexual genesis or HGT, mind in the cosmos, and the perfect continuum. Some scientists believe our cosmos is entangled with another (Chown 2007). On the holographic quantum level, perhaps this is just the competitive entanglement of cI and cro proteins cycling in perpetual motion—an inherent law of nature elucidating the assembly of things. In light of modern experiments, physicist John Wheeler (1988) speculated that the **act of conscious observation functions with quantum mechanics, catalyzing a probability or outcome**. So, if an observing consciousness has knowledge of the classical path and its magnetic field orientation, an act of measurement or choice might entangle space to reify the holographic quantum paths to a perfect continuum. Peirce guesses that to be drawn into a new system, a particle must have the right mass, the right velocity in the right direction, the right attraction, and it must present itself at the right point (EP 1. 270). Similarly, a finite probability exists for any particle that approaches the black hole event horizon to bounce back, dependent on the incoming particle's energy, its charge, and its projection of the orbital momentum on the axis of rotation of the black hole (Kuchiev 2003; 2004; 2004a), the same conditions in Egyptian texts (King 2004; 2006).Only an approximate conclusion remains: chance may be knowledge of the cosmos' holographic mode of operation and magnetic fields; habit breaking may result in the transformation of energy by a violation of the second kind; and evolution of mind or energy might be possible through HGT via a cosmic viral code of Lambda-genesis. "Order is simply thought embodied in arrangement;" (CP 6.490). Accordingly, the original meaning of the signs may relate to a viral biophysics of crystallized mind for human evolution via a natural chemical pathway to a perfect continuum. Again, this semiotics of evolvability and order recurs in time with the chaotic dream of reason that projects a pervasive falsification of our perceptual world. As Wallace Stevens envisions in "Connoisseur of Chaos", perhaps "the pensive man may see." **Peircean Process Metaphysics Origin of Science Quantum and Classical Laws Ancient Egyptian Evolutionary Biophysics** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely

dissimilar are the views of epistemologists in the analysis of raison d'etre of knowledge. Epistemology is the study of knowledge and those things closely related to it: justification, what it takes for you to be justified, the relation between knowledge and justification, whether you can have any justified beliefs at all, and if so, how you come to know (or justifiably believe) things, how you can use what you know (or justifiably believe) to come to know (or justifiably believe) other things, and whether and why it's valuable to know instead of merely having justified and true beliefs. The literature on epistemology is vast. Here's a very brief summary of some epistemological discussions. Concerning knowledge, many epistemologists think knowledge is justified true belief, where the justification you have is linked to the truth of the matter in the right kind of way, though what this way is a matter of debate; some epistemologists think knowledge can't be analyzed this way. Concerning justification's relation to knowledge, some epistemologists think we don't need to be justified to know, and some think we do need to be. Concerning what it takes to be justified, some epistemologists think that what it takes for you to be justified are only factors internal to the believer (Internalists). Others think it takes an external factor, like reliable or well-functioning cognitive faculties (Externalists). Skeptics argue that we can't have any justified beliefs at all, and many epistemologists reply to the skeptic's arguments. Concerning how we use what we know (justifiably believe) to come to know (justifiably believe) other things, some epistemologists (Foundationalists) argue that there are bedrock propositions that we know (justifiably believe), and we build our knowledge (justified beliefs) on those. Others (Coherentists) argue that there aren't bedrock propositions; rather, a set of beliefs is justified as a whole, and several beliefs can be mutually supporting. Concerning the value of knowledge, some argue that knowledge is intrinsically valuable. Others have argued that knowledge is valuable only because of the role it plays in practical reasoning, and others argue that knowledge isn't more valuable than justified and true belief, but there are other epistemic states such as understanding, that do have value above their proper subparts. **Edited y Matthew McGrath (University of Missouri, Columbia). It is not as if there is an anathema or misnomer in the philosophy or physics. Infact**

**most of the pnysicists have used philosophical ideas towards the end of consummation of their their and their concomitant justification thereof.** The eternal return produces becoming-active. It is sufficient to relate the will to nothingness to the eternal return in order to realize that reactive forces do not return. However far they go, however deep the becoming-reactive of forces, reactive forces will not return. The small, petty, reactive man will not return. Affirmation alone returns, this that can be affirmed alone returns, joy alone returns. Everything that can be denied, everything that is negation, is expelled due to the very movement of the eternal return. We were entitled to dread that the combinations of nihilism and reactivity would eternally return too. The eternal return must be compared to a wheel; yet, the movement of the wheel is endowed with centrifugal powers that drive away the entire negative. Because Being imposes itself on becoming, it expels from itself everything that contradicts affirmation, all forms of nihilism and reactivity: bad conscience, ressentiment..., we shall witness them only once. [...] The eternal return is the Repetition, but the Repetition that selects, the Repetition that saves. Here is the marvelous secret of a selective and liberating repetition. There is no need to remind the reader that neither the image of a centrifugal movement nor the concept of a negativity-rejecting repetition appears anywhere in Nietzsche's writings, and indeed Deleuze does not refer to any text in support of this interpretation. Further, one could highlight that Nietzsche never formulates the opposition between active and reactive forces, which constitutes the broader framework of Deleuze's interpretation. For some years, Marco Brusotti has called attention to the fact that Deleuze introduced a dualism that does not exist in Nietzsche's writings. To be sure, the German philosopher describes a certain number of **"reactive" phenomena (for example, in the second essay of the Genealogy of Morality, § 11, he talks about "reactive affects" [reaktive Affekte], "reactive feelings" [reaktive Gefühlen], reactive men [reaktive Menschen]); but these are nonetheless the result of complex ensembles of configurations of centers of forces that remain in themselves active.** Neither the word nor the concept of "reactive forces" ever appears in Nietzsche's philosophy. **After the discovery of the two principles of thermodynamics began a debate about the dissipation of energy and the**

**thermal death of the universe which framed the modern renewal of the debate between the linear and circular conceptions of time**. Scientists such as Thomson, Helmholtz, Clausius, and Boltzmann and-by way of Kant, Hegel and Schopenhauer-philosophers such as Dühring, Hartmann, Engels, Wundt and Nietzsche have tried to address this problem by using the force of scientific argumentation and of philosophical discussion. Whoever believed in an origin and a final end to the motion of the universe (be it in the physical form of the gradual loss of heat, or in the metaphysical form of a final state of the "world process"), relied on the second principle of thermodynamics or on the demonstration of the thesis of Kant's first cosmological antinomy. On the contrary, those who refused to admit a final state to the universe used Schopenhauer's argument of infinity a parte ante-according to which if a final state were possible, it should already have established itself in the infinity of time past-to propose henceforth a number of alternative solutions. Scientists would propose the hypothesis that energy could have re-concentrated after a cosmic conflagration, thus reversing the tendency towards dissipation. Those belonging to the monistic and materialistic tradition relied on the first principle of thermodynamics and on the infinity of matter, space and time, and regarded the universe as an eternal succession of new forms. A certain critical agnosticism was widespread among scientists and philosophers, oftentimes through a reaffirmation of the validity of Kant's antonymic conflict, this movement avoided to take a stand on specifically speculative issues. Other German philosophers, like Otto Caspari, or Johann Carl Friedrich Zöllner, had reintroduced an organicist and pan-psychical conception of the universe, investing atoms with the ability to escape any state of balance. Indeed, it is probably one of Otto Caspari's works, The Correlation of Things (Der Zusammenhang der Dinge. Gesammelte philosophische Aufsätze (Bleslau: Trewendt, 1881)), which awakened Nietzsche's interest for all things cosmological, in that summer of 1881, in Sils-Maria. addressing Schopenhauer and Eduard von Hartmann's mystical pessimism according to which the world is the creation of a stupid and blind essence (which, after having created the world by mistake, comes to the realization that it had made a mistake and strives to return it to nothingness) Caspari stresses that it is

nothing short of mystical to imagine that the world may have been borne out of a an originary and undifferentiated state. Where would it have drawn the first impulse? But, continues Caspari, even if the world had received this first impulse from some deus ex machina, there is no doubt that, in the temporal infinity of past time thus far, it would have either attained the end of the process (but this is impossible because the world would then have ended), or it would be necessarily bound to repeat indefinitely this original mistake, and the entire process that accompanies it. But then, what is the process of the world? We must now take one more step back and understand further the process of the world according to von Hartmann. Hartmann objects that the regressive movement postulated by Schopenhauer is possible only in thought: it remains nothing more than an "ideal postulate" with no real object and which "does not teach us anything about the real process of the world that unfolds in a movement contrary to this backwards movement of thought" (Hartmann, Philosophie des Unbewussten, third edition (1871): 772). Hartmann affirms that if unlike Schopenhauer one admits the reality of time and of the world process, one must also admit that the process must be limited in the past and therefore that there must be an absolute beginning. In Hartmann's mind, failure to do so would result in positing the contradictory **concept of an accomplished infinity**: "The infinity that from the point of view of regressive thinking, remains an ideal postulate, which no reality may correspond to, must, for the world, whose process is, on the contrary, a progressive movement, open up to a determinate result; and here the contradiction comes to light" (Hartmann (1871): 772). What really "comes to light" in this passage is the fact that Hartmann does not provide a demonstration but a petitio principii. Indeed, the concept of the world process analytically contains the concept of a beginning of the world. In all rigor, it is therefore impossible to demonstrate these concepts with reference to each other. Secondly, Hartmann's view that one is bound to accept the reality of the world process even if one rejects the ideality of Schopenhauer's time is mistaken. Hartmann believes that if time is real there must be a world process with both an **absolute beginning and an absolute end. Without any justification, Hartmann jumps from Schopenhauer's negated time to oriented time.** Hartmann's view is that the

world process leads into a final state absolutely identical to the initial state. However, it follows from this that even as the cosmic adventures of the unconscious come to a close, we are still haunted by the specter of a new will and of another beginning of the world process. This exposes a serious internal flaw of Hartmann's system insofar as it jeopardizes the possibility of a final liberation from existence and suffering. This is why in the last pages of his work, "On the Last Principles," he painstakingly calculates the degree of probability of a reawakening of the volitional faculty of the unconscious. Insofar as the will is entirely free, unconditioned and a-temporal, the possibility of a new volition is left to pure mathematical chance and is therefore ½. Hartmann further stresses that if the will were embedded in time, the probability of the repetition would amount to 1 and the process of the world would be bound to begin again, in an eternal return which would completely preclude the possibility of a final liberation. Fortunately, this is not the case since-according to Hartmann's remarkable logic-the world-process develops through time, but the original will is outside of time. In fact, one may even affirm, along the lines of Hartmann's peculiar theory of probability, that every new beginning gradually reduces the probability of the next beginning: let n be the number of times that the will is realized, the probability of any new realization is ½ n. "But it is clear that the probability ½ n diminishes as n increases, in a way that suffices to reassure us in practice" (Hartmann (1869): 663). In his famous work entitled **On the Conservation of Force (1847), Hermann von Helmholtz** had divided the totality of the energy in the universe between potential energy and kinetic energy and affirmed the reciprocal convertibility of the two. In 1852, William Thomson pointed out that there exists a sub-ensemble within kinetic energy, heat, which, once it has been generated, is no longer entirely convertible into potential energy-or into any other form of kinetic energy. Considering that the (partial) reconversion of heat into labor is possible only in situations that present a disparity in temperature, and that heat tends to pass from warmer to cooler bodies by spreading on an even temperature level through space, Thomson concluded that the universe tends towards a final state where any energetic transformations, every movement and every form of life will cease.

Caspari's atoms (which bring to mind those in Leibniz's Monadology) resemble some sort of biological monads, endowed with internal states. For Caspari, every atom obeys the ethical imperative to participate in the conservation of the general organism and its movement does not only follow the simple physical kind of interaction but also an a priori law ensuring. If a balance of forces had been attained at any moment, this moment would still be going on: therefore, it never happened. The present state contradicts this proposition. Supposing that a certain state rigorously identical with the present state had, one day, existed, this supposition is not refuted by the present state. As one of the infinite possibilities, it is necessary that the present state had been given anyway, since until now an infinite period of time has already unfolded. **If equilibrium were possible, it must have occurred**; and if the present state has already taken place, then so too the one that preceded it as well as the one preceding that one. Therefore it has already taken place a second time, a third time and so on. And likewise it shall take place again a second time, a third time. Innumerable times forwards and backwards. This amounts to saying that all becoming occurs within a repetition of an innumerable number of absolutely identical states. [.] The immovability of forces, their equilibrium is a conceivable case, but it has not occurred. As a result the number of possibilities is greater than the number of realities. -The fact that nothing identical recurs may be explained not thanks to chance, but only thanks to an intention infiltrated within the essence of force. Indeed, supposing an enormous amount of cases, the random occurrence of the same combination is more probable than the same combination never recurring. The quantum of force in the universe is determinate and not "infinite": let us beware [hüten wir uns] from such conceptual extravaganza! Therefore the number of situations, modifications, combinations and developments of this force is doubtless enormous and practically "immeasurable," but in any case this number is determinate and not infinite. On the other hand, the time in which the universe exerts its force is infinite. That is to say, that force is eternally identical and eternally active: -until the present instant an infinity has already taken place, that is to say that all possible developments must have already taken place. Consequently, the present development must be a repetition

and therefore both this that was born from it and this that shall be born from it and so on both forwards and backwards. Everything has taken place an innumerable number of times because the overall situation of all forces always recurs (FP 11 [202] of 1881). Nietzsche regards the mechanistic vision as more plausible and less anthropomorphic than organicism. However, faced with the two major **cosmological models of his time, the mechanistic model and the organic model, Nietzsche wishes to return its polymorphous, proteiform, unstructured** and chaotic character to nature of which the perfectly non-theological and non-teleological theory of eternal return is the strongest seal. This is the first of the "new battles" which come to whoever is aware of the consequences of the death of God: take any antropomorphism away from nature. In the preparatory papers, the third book of the Gay Science is entitled "Gedanke eines Gottlosen / Thoughts of a Godless One." Aphorism 109 of this book, which immediately follows the famous aphorism against the shadows of God, **Boltzmann's theory belongs to the third phase of the debate on thermodynamics and cosmology. The publication of Thomson's brief paper "On a Universal Tendency in Nature to the Dissipation of Mechanical Energy" of 1852 signals the beginning of the scientific controversies on the problem of the dissipation of energy and opened the first phase of the debate, which was announced in the Reflexions on the Motor Powers of Fire by Sadi Carnot and whose conclusion is represented by Clausius's recapitulative article on the concept of entropy in 1865.[44] The second phase started in 1867 when, at the forty-first congress of German scientists and doctors, Clausius gave a lecture on "The Second Principle of the Mechanistic Theory of Heat," where he applied the results of his research on thermodynamics to the universe**. It is true that in his famous lecture of 1854, Helmholtz had already presented the cosmic consequences of the second principle, but Clausius' contribution had a strong impact on German culture. This is because in this lecture he robustly rejected the possibility to consider the universe as an eternal and self-renewing circle, an ewiger Kreislauf in which force and matter are in constant transformation, as was heretofore affirmed by the materialism of the scientists and philosophers, and he did so in the name of the

second principle of thermodynamics. In this way, the debate on the principles of thermodynamics gained great importance in European Culture starting in 1867. In the two first phases, it is Thomson's mechanism that predicts the thermal death of the universe. In the third phase, on the contrary, the meaning of the term mechanism changes radically.[46] In accordance with the apocalyptic climate of this period dominated by the "rebirth of idealism," the "overcoming of scientific materialism" and the "bankruptcy of science", the mechanistic paradigm which had accompanied the birth of modern science became challenged on **account of the second principle of thermodynamics**. According to the theorem of the quasi-periodicity of the motions of mechanical systems demonstrated by Poincaré as part of the problem of the three bodies (1890), a mechanical system must evolve according to a quasi-periodical movement and consequently it must always return-sooner or later-to the initial state. **An easily established theorem informs** us that a limited world obeying solely the laws of mechanics shall always pass through a state closely similar to its initial state. On the contrary, according to established experimental laws (supposing we grant them absolute value and wish to push their consequences to the end), the universe is directed towards a final state, which once it is attained, it shall not be able to escape. In this final state, which shall be like a sort of death, all material bodies shall be at rest at the same temperature. **Poincaré's theorem seems** therefore incompatible with the second principle of thermodynamics, which predicts a unidirectional movement of all natural phenomena until the whole universe is brought to a total standstill. Wilhelm Ostwald and the entire energeticist school of thought contended that the principles of thermodynamics were fundamentally new, and could not be re-incorporated to traditional physics and that they should serve as a basis for a new science that regards the qualitative diversity of energy and its tendency to degradation as its axioms. Against energeticism and in an effort to bring entropic phenomena back into the theoretical framework of mechanism, Ludwig Boltzmann introduced the concept of probability in physics, not as an instrument of calculation, but as an explicative principle. In Boltzmann's statistical thermodynamics, the increase of entropy assumed by Clausius is re-interpreted as an increase in molecular

chaos. As a result, it becomes possible to explain mechanistically the evolution of closed systems endowed with increasing entropic value, without it committing us to granting absolute value to the second principle of thermodynamics. **Moreover, one no longer needs to fear the thermal death of the universe insofar as the state of equilibrium will in principle** never be complete, but rather will be attained only statistically, leaving open the possibility of fluctuations towards less probable states. Boltzmann's critics remarked that this hypothesis involved two paradoxes called the objection of reversibility (Umkehreinwand), and that of repetition (Wiederkehreinwand). I shall only address here the second one since it coincides with the theory of the eternal return. Based on Poincaré's theorem quoted above, Ernst Zermelo objected to Boltzmann that his model of the universe suggested that after a finite (if admittedly very long) time the system would return to its initial position. In his first response to Zermelo, Boltzmann avoids committing himself directly to cosmological questions and he only observes that, in the case of concrete thermodynamic systems, the time of recurrence may be extremely long. For example, in normal conditions of pressure and temperature, one-centimeter cube of gas requires 1010 years to reach a molecular configuration identical to the original one! However, following a response by Zermelo, Boltzmann wrote a new article where he outlines a cosmological picture that he will re-use later in his conclusion to his famous Lectures on Gas Theory. In this cosmological picture, Boltzmann considers the universe as a closed system with constant entropy, within which some fluctuations occur, creating islands of negative entropy. Our solar system originates in one of these fluctuations. As Clausius correctly pointed out, the entropy of our solar system increases constantly as the solar system gets closer to the state of chaos and of the thermal death of the rest of the cosmos. However, in other zones of the universe, some new fluctuations and new islands appear, so that thermal death is never generalized. Here we are given a grand cosmic image, in which the solar system and the sparkle of life that was lit on planet earth are only a fluctuation of order from within a dominant entropic tendency. Life, and the order on which it is based are exceptions, transitory forms taking place in the realm of the shapeless, they are islands of the cosmos that will soon

be re-absorbed into chaos. According to Poincaré's theorem, our island will have to be reborn, to develop and die innumerable times in a strictly identical fashion. This happened an infinite number of times during the past eternity and it shall take place again an infinite number of times in the eternity to come. **Boltzmann accepts the "paradox" of recurrence-that is the eternal return of the same-as a legitimate consequence of the probabilistic conception of thermodynamics**. It may be rejected for ethical reasons, it may be stored away as an abstract speculation or dismissed along with other cosmic fantasies, but it cannot be rejected on the basis of any rigorously scientific viewpoint. **The Eternal Return: Genesis and Interpretation BY PAUL D'IORIO(excerpts). Nietzsche found the template of a materialistic cosmology based upon the first principle of thermodynamics. During the same year, he could have found a model of an organistic solution to the problem of thermal death of the universe as well as a discussion on the conformation of space in Zöllner's book On the Nature of Comets** (Johann Carl Friedrich Zöllner, Über die Natur der Kometen. Beiträge zur Geschichte und Theorie der Erkenntnis (Leipzig, Engelmann, 1872), 299 f. and 313 f.); Assumed Association In Space Time Of Particles Or Men Leads To Detrimental Ramification Be In Terms Of Progress Of “Subject Matter” Or In The Pathological Disequilibrium Of Gratification Deprivation Complex Which Again Has Mental Turbulence In Case Collective Consciousness Approve It (Raaga) Or Pernicious Implications In Case The Theory Is Not Accepted; After All What Is A Theory: Set Of Correlations And Causations Which Are “Tested” Against A Collective Consciousness Background, Which Yields “Results’ Which Are Themselves In “Question” Due To The Very Perception And The Simulation Prospective. The Definition Of Quantum Theorists' Terms, Such As Wavefunctions And Matrix Mechanics, Progressed Through Many Stages. For Instance, Erwin Schrödinger Originally Viewed The Electron's Wavefunction As Its Charge Density Smeared Across The Field, Whereas Max Born Reinterpreted It As The Electron's Probability Density Distributed Across The Field. Although The Copenhagen Interpretation Was Originally Most Popular, Quantum Decoherence Has Gained Popularity. Thus The Many-Worlds Interpretation Has Been Gaining

Acceptance Moreover, The Strictly Formalist Position, Shunning Interpretation, Has Been Challenged By Proposals For Falsifiable Experiments That Might One Day Distinguish Among Interpretations, As By Measuring An AI Consciousness Or Via Quantum Computing. More Or Less, All Interpretations Of Quantum Mechanics Share Two Qualities: They Interpret Formalism—A Set Of Equations And Principles To Generate Predictions Via Input Of Initial Conditions They Interpret A Phenomenology—A Set Of Observations, Including Those Obtained By Empirical Research And Those Obtained Informally, Such As Humans' Experience Of An Unequivocal World Two Qualities Vary Among Interpretations: Ontology—Claims About What Things, Such As Categories And Entities, Exist In The World Epistemology—Claims About The Possibility, Scope, And Means Toward Relevant Knowledge Of The World. Now How Does A Person In A Crowd Behave? They He Would Due To Herd Mentality; Ok; That Is One Theory; Simulation By Terrestial Brahman Anti Brahman Will Make Them Behave In A Highly Organised And Coordinated Manner; With Association Of Things And Individuals Already Established Resultant Orientation Of The Actions Hall Be Only Oriented Towards The Figment Of Imagination And Product Of Puerile Prognostication Formally Established By Terrestial Brahman Anti Brahman Before Such A Coordinated And Organised Execution Of Exercise Is Conducted; Analogically, If A Simulation Is Going On In A Bunch Of Particles Or Electrons, And If Some Agency Say Celestial Brahman Anti Brahman Agency Is Simulating It, There Shall Be Coordination And Organisation So That You See “Reality” While The Truth Shall Remain Hidden. Result Is The Suppression Of Information, Delineation Of Disinformation, And The Concomitant Development Of Mushrooms Of Concocted And Fabricated Theories To That “Reality”; In The Case Of Terrestial Brahman Anti Brahman Agency Intention May Be To Bulldoze And Bowdlerize You, Baffle And Befuddle You In To Submission; In The Case Of Celestial Brahman Anti Brahman Agency Only Reason Could Be To Shield The Truth Being Disseminated; But Why?;May Be You Should Never Become God Which You Are Trying To; May Be God Is Really Playing Dice; What Does That Result In To ; You Become A Scezophreniac Associating Again And Again What You Had Been

Until Now Of Particles And Other Concomitants Again And Again Leading To Flawed Inferences. Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. In A Broad Sense, Scientific Theory Can Be Viewed As (=) Offering Scientific Realism—Approximately True Description Or Explanation Of The Natural World—Or Might Be Perceived With Antirealism. A Realist Stance Seeks The Epistemic And The Ontic, Whereas An Antirealist Stance Seeks Epistemic But Not The Ontic. In The 20th Century's First Half, Antirealism Was Mainly Logical Positivism, Which Sought To Exclude Unobservable Aspects Of Reality From Scientific Theory. Since The 1950s, Antirealism Is More Modest, Usually Instrumentalism, Permitting Talk Of Unobservable Aspects, But Ultimately Discarding The Very Question Of Realism And Posing Scientific Theory As A Tool To Help Humans Make Predictions, Not To Attain Metaphysical Understanding Of The World. The Instrumentalist View Is Carried By The Famous Quote Of David Mermin, "Shut Up And Calculate", Often Misattributed To Richard Feynman. Quantum Mechanical Wave Function (Absolutely Squared) Describes The Completed Interference Pattern, It Must Describe An Ensemble Interpretations Of Quantum Mechanics From Wikipedia. Philosophy Of Science, The Distinction Of Knowledge Versus Reality Is Termed Epistemic Versus Ontic. A General Law Is A Regularity Of Outcomes (Epistemic), Whereas A Causal Mechanism May Regulate The Outcomes (Ontic). A Phenomenon Can Receive Interpretation Either Ontic Or Epistemic. For Instance, Indeterminism May Be Attributed To Limitations Of Human Observation And Perception (Epistemic), Or May Be Explained As A Real Existing Maybe Encoded In The Universe (Ontic). Confusing The Epistemic With The Ontic, Like If One Were To Presume That A General Law Actually "Governs" Outcomes—And That The Statement Of A Regularity Has The Role Of A Causal Mechanism—Is A Category Mistake. Interpretations Of Quantum Mechanics From Wikipedia. Reality Is Constant In The Midst Of Change. Note The Conservation Of Individual, Collective And Cosmic General Ledgers Are What Are Being Spoken Of. Not The Changing

Transactions, But The Immutable Refractory Zero. (Italics Mine) What This Means Essentially Is That There Is Change Although Nothing Changes. Like There Are Changing Transactions And Still The Assets And Liabilities Will Be Equal. This Impossible Situation Is Reflected (E&Eb) In The Ultimate Impossibility Of Change Itself. That Which Does Not Exist Prior To Its Changing And At The End, After It Has Changed, Must Be Equally Non-Existent Between These Two Moments. Although The World Of Change Appears To Be Real, It Cannot Be So. (Real)Change, According To The Vedantin, Presupposes A Loss Of Identity. Reality Cannot Suffer Transformation If It Were To Do So, It Would Become Something Else And The Real Would Be Deprived Of Its Reality. The Immortal Can Never Become Mortal, Nor Can The Mortal Become Immortal The Doctrine Of Vibration: An Analysis Of The Doctrines And Practices Of Kashmir Shaivism Mark S. G. Dyczkowski

**ABSOLUTE STABILITY**

$$L_9 = \frac{N_9 G_9}{2M_9} \frac{d(9nF_9)}{dR_9} (\geq 0) \quad (9a)$$

$$I_9 = -\frac{N_9 G_9}{2M_9} \frac{d(9nM_9}{dR_9} (\geq 0) \quad (9b)$$

It is shown in appendix 2 how the condition (6), which can now be expressed in the form, $P_9 > L_9 - I_9$, must be modified ahe $dM_9/dR_9$ becomes too large. This allows us to determine the necessary and sufficient condition for absolute stability as

$P_9 > L_9 - I_9$, when $P_9^2/4 > I$ (8a)

$P_9 > L_9 - I_9 + \frac{(I_9 - P_9^2}{(L_9 - I)}$ when $P_9^2/4 < I.$ (8b)

The conditions (8) show that the total debit regulation, P, which is not directly affected by the delay, continues to promote stability as it would in the undelayed model$[c.f.condition(2)]$. The delayed term, L, acts directly to

reduce stability as we might expect. Less expected is the role the undelayed component of debit dynamics, $I$, which acts to offset the destabilizing influence of the time delay. However, the immediate effect, $I$,

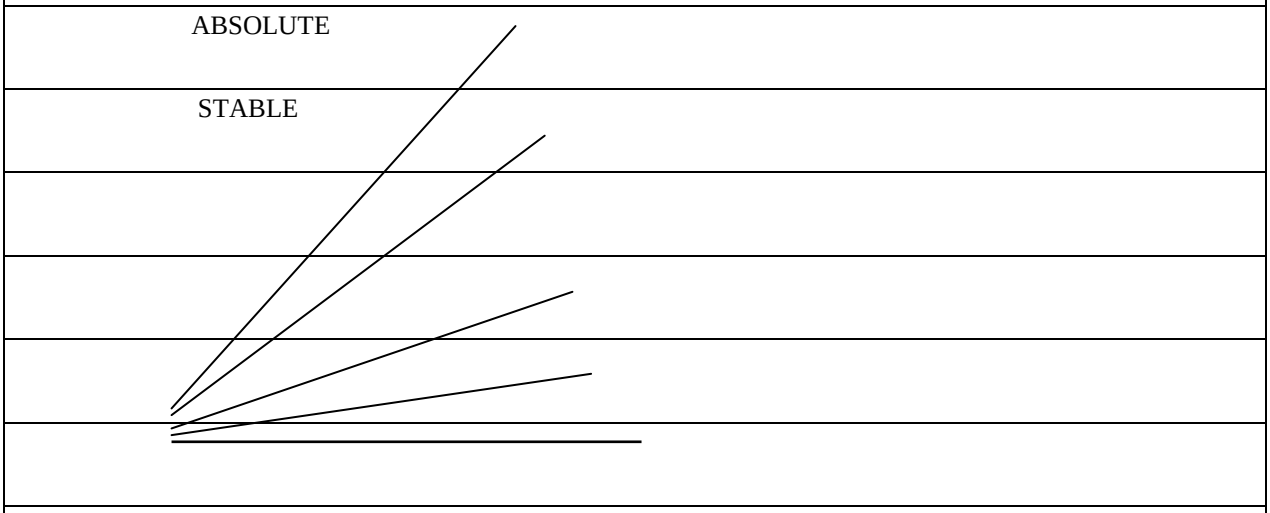



FIG.4.Absolute stability in the presence of a debit delay. In each case the zone of stability is above and to the left of the lines shown. The total debit regulation $(P = (dD/dR + NdG/dR - dB/dR)/M)$ must be positive for an undelayed system to be stable. The delayed credit effect $L = (NGdF/dR)/2M^2)$ reflects the influence of the debit (like in cheque in abbey axe) delay. The different lines show the zone of absolute stability given different proportional relationship between $L$ and the immediate effect $(I = (NGdM/dR)/2M^2)$.

Can never completely offset the lagged effect,$L_9$; condition (8b) ensures that the total debit regulation, $P_9$, must always be positive even when $I_9$ is very large. This pattern is shown in Fig. 4, where the region of absolute stability is illustrated for various proportional relationships between $I_9$ and $L_9$. As $I_9$ increases relative to $L_9$ the zone of stability is enlarged; however,even when $I_9$is zero a significant zone of stability exists.

It is important to note that in considering the parameter space shown in Fig. 4 we need not consider changes in equilibrium values of the predator and prey. These values are determined directly by the functions involved in the dynamics, whereas regulation is determined primarily by the first derivative of these functions. There is no a priori

reason why these two feature should covary and as a result it is possible to traverse all of the parameter space shown in Fig. 4 without affecting the equilibrium values.

As in the case of a prey delay, one reason why the conditions for absolute stability are important is that very short delays can dramatically change stability; indeed this effect is more marked for a credit than a debit delay and is shown in Fig. 5. The two examples illustrated differ in the ratio of $L$ to $I$ and consequently differ considerably in the shape of their stability curves; however, in both a delay of only one- tenth of the natural period results in a dramatic loss of stability. We see that when $I_9$ is relatively large, as in Fig. 5a stability does not always decline in the presence of a time delay; the addition of a time delay of around half the natural period makes the system more stable.

The derivation of the conditions (8), defining absolute stability in the

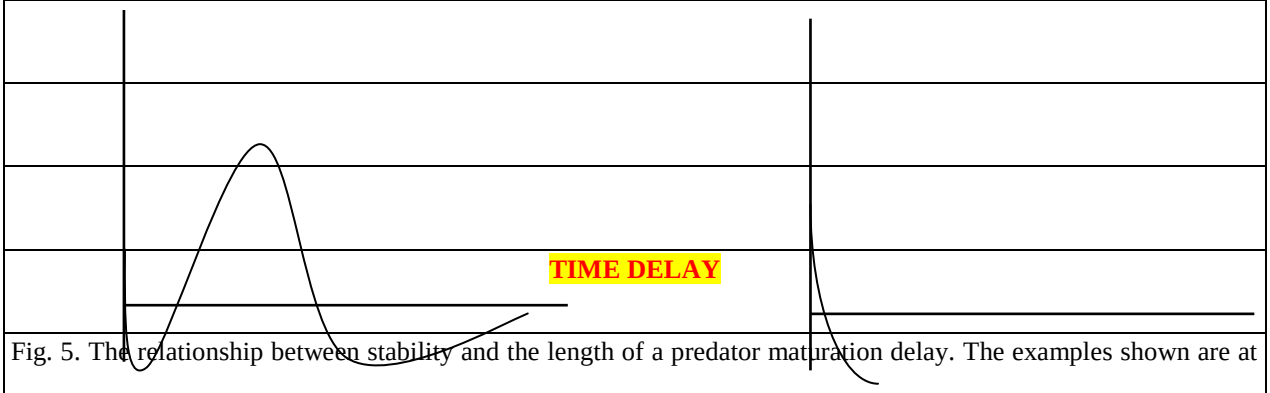



Fig. 5. The relationship between stability and the length of a predator maturation delay. The examples shown are at the limit of absolute stability. (a) The immediate predator effect, I, is large relative to the delayed predator effect, L. (b) The immediate predator effect, I, is small relative to the delayed credit effect, L. I and L are defined as in Fig. 4 and text Eq. (9). Parameters in each case M=9; $dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9 = 0.05$ For (a) $NGdF/dR = 0.3, -N_9G_9dM_9/dR_9 = 0.6$; for (b)$NG_9dM_9/DR_9 = 0$

**DELAY IN BOTH PREDATOR AND PREY VIS A VIS DEBIT TYPES AND CREDIT TYPES**

If both the credits and the debit are subject to a maturation period of T then the model (1) becomes

$$\frac{dN_9t}{dt} = N_9(t - T_9)F_9\big(R_9(t - T_9)\big) - N_9(t)M_9\big(R_9(t)\big)$$

(9)

$$\frac{dR_9t}{dt} = B_9\big(R_9(t - T_9)\big) - D_9\big(R_9(t)\big) - N_9(t)G_9\big(R_9(t)\big)$$

As in the two previous delay models, (3) and (5), it is assumed that the nascent have no impact upon the dynamics of the dynamics of the system except via their effect in delaying the response of the populations to changed conditions.

The condition (4) requires that in the presence of a debit delay both the total prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) and the effective prey regulation ($dD_9/dR_9 + NdG_9/dR_9 - dB_9/dR_9$) must be positive. The condition (8) requires that in the presence of a predator delay the total prey regulation must be greater than some positive value. The size of this value rises as the effect of delayed credit recruitment is increased, although this rise can be offset by a similar increase in the effect of undelayed credit mortality. Combining the general principles of (4) and (8) suggests that when both the credits and the debits are subject to a delay the total debits regulation will be constrained above zero, but not as severely constrained as would be expected from (8). On the other hand we could expect the effective debit regulation to be more severely constrained, reflecting the destabilizing effects of delayed regulation acting through both debit and credit recruitment.

Before examining these expectations we need to realize that the behavior

Of the double – delayed system (9) depends upon the relative importance of the two delays. For example, $dB_9/dR_9$ is a measure of the control acting through the debit delay and setting $dB_9/dR_9 = 0$ removes the dynamic effect of the debit delay. Similarly $dF_9/dR_9$ is a measure of the control acting through the credit delay although setting

$dF_9/dR_9 = 0$ does not negate the effect of a predator delay; the value of $M_9$ which measures the equilibrium per capita creation and notification rates of the credit types, also acts as a vehicle through which the influence of the delay can be felt. In order to investigate the joint effect of a credit and a debit delay we need to use this information and define a credit-debit system which approximately balances the two delays. To this end we set the effect of the debit delay, as measured by the magnitude of $dF_9/dR_9$and normalized by $M_9$ (see9a). The absolute stability of such a system is summarized in Fig. 9 and 8. As was the case for the previous models, it is important to realize that the parameter space shown in these figures can be traversed without altering the equilibrium values. For different values of credits and debits the graph.

A comparison of Fig. 1 with Fig. 4 illustrates that the constraints upon the total debit regulation are markedly decreased by adding a debit delay to a system already including credit delay. In fact, with regard to the total debit regulation, the double-delayed system is hardly more constrained than the undelayed system. Furthermore, this difference becomes

Systems to which model has been applied also include: **Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI:

http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors.

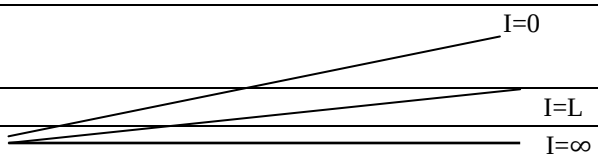
Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, ,**

**Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)**

TOTAL

ABSOLUTELY STABLE

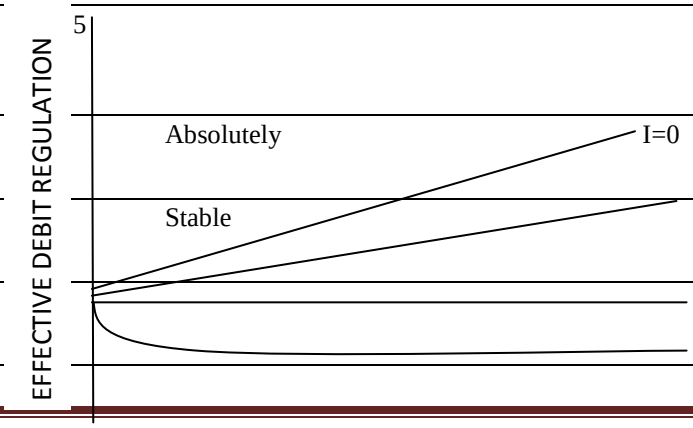



DELAYED DEBIT EFFECT

FIG9. Absolute stability in the presence of debit and credit delays of equal length (and equal importance; see text). When debit recruitment is not regulated ($dB_9/dR_9 > 0$) absolute stability is usually limited by the total debit regulation, P. For further details see Fig.4, which is identically scaled for direct comparison.

Courtesy Nunney : For different values of Haimovician accentuation and attrition coefficies absolutes stability obtain.

However, the effect is limited (see Fig.8) because even when I is infinite the effective prey regulation must be greater than- 2M (see Appendix 3), so that under all conditions there is a danger of overly strong regulation of prey recruitment.



-50 10

**DELAYED CREDIT EFFECT (LACK OF BALANCE) CREDIT NON CUMSUMMATION:"**

Fig.8 Absolute stability in the presence of debit-creditor delay of equal length (and equal importance: see text) When prey recruitment is regulated ($dB/dR < 0$) absolute stability is limited only by the effective prey regulation ($= (dD_9/dR_9 + N_9 dG_9/dR_9 + dB_9/dR_9)/M_9$). The immediate ($I_9$) and lagged ($L_9$) predator effect are defined as for Fig.4 and text Eq.(1). The effective debit regulation is not directly limited in the absences of a delay, but must be positive for absolute stability in the presence of a debit delay .

Vanishingly small as the undelayed part of the debit dynamics, as measured by $I_9$(see 9b), is increased relative to the lagged part, $L_9$

As predicted earlier the pattern is very different when we look at constrains upon the effective debit regulation. Such constraints reflect the fact that when debit recruitment is subject to a delay, strong regulation of the process can be destabilizing. Reference to Fig. 8 shows that the addition of a debtor delay has a fairly straight\forward effect. When the delayed predator effect (L) is more –or-less equal to the immediate predator effect (I) then the effective prey regulation is contained as if there was no predator delay. However, if $L_9 > I_9$then the constraint can be markedly increased. In fact, as Nunney (1985a) has noted, whenever $L_9$is somewhat large than $I_9$ then the condition constraining prey regulation of recruitment ($dB_9/dR_9 < 0$) is

$$\frac{dD_9}{dR_9} + N_9\frac{dG_9}{dR_9} + \frac{dB_9}{dR_9} > \frac{(L_9 - I_9)M_9}{2} \qquad (10)$$

If $L_9 < I_9$ then the constraint upon the left-hand side of (10), the effective debit regulation, is lessened. In fact a large value of $I$ can offset the effect of the credit delay plus part of the effect of the debit delay so that the constraint can become less than when only the debit are subject to a delay.

The following treatise bears ample testimony and infallible observatory to the fact that how Nunny's model can be applied to various cosmological, general relativity and economic problems.

## ONTOLOGICAL AND EPISTEMOLOGICAL FEATURES OF THE STUDY: OF NUNNY's MODEL

When they are known to be as they are, in the fullest sense of their existence, their phenomenal nature disappears leaving the ground of Being naked and accessible. This approach was validated by a critique of experience. The Vedanta established that space, time and the other primary categories of our daily experience can have no absolute existence. It was therefore necessary to make a distinction between relative truths — that accepted by the precritical common man — and an absolute truth discovered at a higher level of consciousness. The Saiva absolutist 17 rejects any theory that maintains that the universe is less than real. From his point of view a doctrine of two truths, one absolute and the other relative, endangers the very foundation of monism. The Kashmiri Saiva approach is integral: everything is given a place in the economy of the whole. It is equally wrong to say that reality is either one or diverse. Those who do so fail to grasp the true nature of things which is neither as well as both. "We do not" says Abhinavagupta, "base our contention that [reality] is one because of the contradictions inherent in saying that it is dual. It is your approach {paksa) that accepts this [method] While, if [duality and oneness] were in fact [to contradict each other], they would clearly be two [distinct realities]." The Vedantin, who maintains that non-duality is the true nature of the absolute by rejecting duality as only provisionally real, is ultimately landed in a dualism between the real and illusory by the foolishness of his own excessive sophistry (vacafadurvidya). Oneness is better understood as the coextensive unity (ekarasa) of both duality and unity. They are equally expressions of the

absolute. Gopinath Kaviraj says: According to Sarikara, Brahman is truth and Maya is inexplicable (anirvacaniya). Hence the [Advaitin's] endeavour to demonstrate the superiority of Advaita philosophy is turned against his own system. It tarnishes the picture of its philosophical perfection and profundity. He cannot accept Maya to be a reality, therefore his non-dualism is exclusive. The whole system is based on renunciation and elimination and thus is not all-embracing.... By accepting Maya to be Brahman (brahmamayi), eternal (nitya) and real (satyarupa), Brahman and Maya [in the Tantra] become one and coextensive. **The Doctrine of Vibration: An Analysis of the Doctrines and Practices of Kashmir Shaivism Mark S. G. Dyczkowski. Concept of being from Shaivite point of view:** Solutional behaviour and stability analysis of such systems is of paramount importance, notwithstanding the assumptions made which are not farfetched, in that it throws light on the fact whether they could be candidates for yet unsolved problems, providing Autheticative determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication for the wide amplitudinal ramificatory usage of the thesis propounded. The world is now the other person's world, a foreign world that no longer comes from the self, but from the other. The other person is a "threat to the order and arrangement of your whole world...Your world is suddenly haunted by the Other's values, over which you have no control"(Satre: being and Nothingness) Sartre states that many relationships are created by people's attraction not to another person, but rather how that person makes them feel about themselves by how they look at them. This is a state of emotional alienation whereby a person avoids experiencing **their subjectivity by identifying themselves with "the look" of the other**. The consequence is **conflict**. In order to maintain the person's own being, the person must control the other, but must also control the **freedom of the other "as freedom**". These relationships are a profound manifestation of "bad faith" as for-itself is replaced with the other's freedom. The purpose of either participant is not to exist, but to maintain the other participant's looking at them. This system is often mistakenly called "love", but it is, in fact, nothing more than emotional alienation and denial of freedom. In

essence we are always simulating and selling ourselves at the cost of our freedom. Both as a script writer and participant consciousness, human beings experienceâ complete subjective feelings and sometimes absolute subjective feelings, which if remains unchecked might lead to pathological ramifications. (Italics mine) Perspicuous forbearance, Sophisticated seasoning, Partcipational observation, Background actuality, Existential worldliness, Existential strife, Predicational anteriorities, Character consonance, Ontological consonance, Primordial exactitude, Phenomenological correlates, Accolytish representation, Atrophied asseveration. Anamensial alienisms, Anchorite aperitif, Arcadian Atticism all delineated in the conditionalities and functionalities and orinetatilities of the system which incorporate the rules and regulations, axiomatic predications and postulation alcovishness of the foregoing state form the bastion, pillar, post, stylobate and sentinel of the classification scheme and doxa. appropriate and diagnostic, differentiating-discriminating, distinctive-distinguishing, emblematic- especial, essentially exclusive, fixed, idiosyncratic-inborn, inbred, indicative and individual, individualistic-inherent, individualizing and ingrained, inherent, innate, local, marked, native, normal, original, particularly- peculiar, personal, private, proper, regular, representative-singular, special, specific, symbolic, symptomatic- unique these form the classification and stratification measure that is suitable notwithstanding the constancy in many a case, as in gravity, which depends on the mass and distance for two given objects**.** Like Kant and Bergson, Deleuze considers traditional notions of space and time as unifying forms imposed by the subject. He therefore concludes that pure difference is non-spatio-temporal; it is an idea, what Deleuze calls "the virtual". (The coinage refers to Proust's definition of what is constant in both the past and the present: "real without being actual, ideal without being abstract.") Proust, Le Temps Retrouvé, ch. III: see the fourth line from the bottom of this page, or, in English translation, the thirteenth paragraph here: "I began to discover the cause by comparing those varying happy impressions which had the common quality of being felt simultaneously at the actual moment and at a distance in time, because of which common quality the noise of the spoon upon the plate, the unevenness of the paving-stones, the taste of the madeleine, imposed the past

upon the present and made me hesitate as to which time I was existing in. Of a truth, the being within me which sensed this impression, sensed what it had in common in former days and now, sensed its extra-temporal character, a being which only appeared when through the medium of the identity of present and past, it found itself in the only setting in which it could exist and enjoy the essence of things, that is, outside Time. [...] **Nothing but a moment of the past Much more perhaps; something which being common to the past and the present, is more essential than both. [...] a marvellous expedient of nature had caused a sensation to flash to me—sound of a spoon and of a hammer, uneven paving-stones**—simultaneously in the past which permitted my imagination to grasp it and in the present in which the shock to my senses caused by the noise had effected a contact between the dreams of the imagination and that of which they are habitually deprived, namely, the idea of existence—and thanks to that stratagem had permitted that being within me to secure, to isolate and to render static for the duration of a lightning flash that which it can never wholly grasp, a fraction of Time in its pure essence. When, with such a shudder of happiness, I heard the sound common, at once, to the spoon touching the plate, to the hammer striking the wheel, to the unevenness of the paving-stones in the courtyard of the Guermantes' mansion and the Baptistry of St. Mark's, it was because that being within me can only be nourished on the essence of things and finds in them alone its subsistence and its delight. It languishes in the observation by the senses of the present sterilised by the intelligence awaiting a future constructed by the will out of fragments of the past and the present from which it removes still more reality, keeping that only which serves the narrow human aim of utilitarian purposes. But let a sound, a scent already heard and breathed in the past be heard and breathed anew, simultaneously in the present and in the past, real without being actual, ideal without being abstract, then instantly the permanent and characteristic essence hidden in things is freed and our true being which has for long seemed dead but was not so in other ways awakes and revives, thanks to this celestial nourishment "While Deleuze's virtual ideas superficially resemble Plato's forms and Kant's ideas of pure reason, they are not originals or models, nor do they transcend possible experience; instead

they are the conditions of actual experience, the internal difference in itself. "The concept they [the conditions] form is identical to its object."[Desert Islands, p. 36.] A Deleuzean idea or concept of difference is therefore not a wraith-like abstraction of an experienced thing, it is a real system of differential relations that creates actual spaces, times, and sensations. [See "The Method of Dramatization" in Desert Islands, and "Actual and Virtual" in Dialogues. In fact it not as if the style and substance of science and philosophers have been entirely different. look at the conformality and consonance in the Gaussian Theory and Deleuze's statements on identity and the loss of it in the following passage…….That is, identity is a continuous process that finds its limit in exhaustion. As Deleuze will write, "…the logical relation of causality is inseparable from a physical process of signalling, without which it would not be translated into action. By 'signal' we mean a system with orders of disparate size, endowed with elements of dissymmetry; by 'sign' we mean what happens within such a system, what flashes across the intervals when a communication takes place between disparates. The sign is indeed an effect, but an effect with two aspects: in one of these it expresses, qua sign, the productive dissymmetry; in the other it tends to cancel it" (dr, 20). An actualized entity just is such a signal-sign system, remaining in perpetual communication with the virtual multiplicity it actualizes, constituting its own elements and identity. Such is Deleuze's account of what Badiou refers to as the operation of the "count-for-one"… an operation that is immanent to being itself and energetic in character. yet how are we to discern the local nature of ontological situations, the manner in which there is no global, overarching global situation in which all local situations are embedded as if parts of a whole? in part the above passage already responds to this question. if series must be brought into relation to communicate in order for causality to occur, and if the resonance of series is always a matter of chance, of a throw of the dice, like Lacan's objet a where no series enjoys primacy of model over copy, then all we have are local situations without a total situation. Unlike Heidegger's analysis of Dasein where Dasein is always being-in-the-world, Deleuze's simulacra or actualizations are always divergent, enjoying only local relations. Heidegger argues that in order to

understand the being of Dasein we must understand the manner in which Dasein is a being-in-the-world. By contrast, Deleuze's concept of multiplicity provides us with the principle of a new structuralism, of a new local ontology, that allows us to understand the immanent organization of a multiplicity without referring it to an embedding global space. as DeLanda so beautifully puts it in relation to the mathematician gauss, "…when gauss began to tap into these differential resources, a curved two-dimension surface was studied using the old Cartesian method: the surface was embedded in three-dimensional space complete with its own fixed set of axes; then, using those axes, coordinates would be assigned to every point of the surface; finally, the geometric links between points determining the form of the surface would be expressed as algebraic relations between the numbers. but gauss realized that the calculus, focusing as it does on infinitesimal points on the surface itself (that is, operating entirely with local information), allowed the study of the surface without any reference to a global embedding space. basically, gauss developed a method to implant the coordinate axes on the surface itself (that is, a method to implant the coordinate axes on the surface itself (that is, a method of 'coordinatizing' the surface) and, once points had been so translated into numbers, to use differential (not algebraic) equations to characterize their relations. as the mathematician and historian Morris Kline observes, by getting rid of the global embedding space and dealing with the surface through its own local properties, 'gauss advanced the totally new concept that a surface is a space in itself'" (intensive science and virtual philosophy, 11-12).according to the old Cartesian method, we can only outline the properties of a space by relativizing it to a global space in terms of which it is then mapped. by contrast, gauss is able to explore a space in terms of its intrinsic metric and organization as a local space, without referencing it to a whole of which it is conceptualized as a part. **The Image of Thought** "Image of thought "permeates both popular and philosophical discourse. According to this image, thinking naturally gravitates towards truth. Thought is divided easily into categories of **truth and error**. The model for thought comes from the educational institution, in which a master sets a problem and the pupil produces a solution which is either true or false. This image of the subject

supposes that there are different faculties, each of which ideally grasps the particular domain of reality to which it is most suited. In philosophy, this conception results in discourses predicated on the argument that "**Everybody knows..."** the truth of some basic idea. Descartes, for example, appeals to the idea that everyone can at least think and therefore exists. Deleuze points out that philosophy of this type attempts to eliminate all objective presuppositions while maintaining subjective ones. Deleuze maintains, with Artaud, that real thinking is one of the most difficult challenges there is. Thinking requires a confrontation with stupidity, the state of being formlessly human without engaging any real problems. One discovers that the real path to truth is through the production of sense: the creation of a texture for thought that relates it to its object. Sense is the membrane that relates thought to its other. Accordingly, learning is not the memorization of facts but the coordination of thought with a reality. "**As a result, 'learning' always takes place in and through the unconscious, thereby establishing the bond of a profound complicity between nature and mind".** Deleuze's alternate image of thought is based on difference, which creates a dynamism that traverses individual faculties and conceptions. This thought is fundamentally energetic and signifying: if it produces propositions, these are wholly secondary to its development. Philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; even tum tan tum;it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the vorticity of being merges with the positive use of the disjunctive synthesis, and

this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoin der for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for every hum of voices or a single silence for all the silences; a single vacuum for all the vacuousness consciousness should not be said without conjuring if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. The attributable to all states of bodies and states of affairs and the expressible of every proposition So univocity of consciousness means the identity of thé noematic attribute and that which is expressed linguistically and sense-fully. Univovity.We mean that it does not allow consciousness to be subsisting in aquasi state and but expresses in all pervading reality. (Reference: LOS: Deleuze) Deleuze sometimes spoke of getting up behind an author and "creating a monster", as if his reading methodology somehow consisted of a distortion of the philosopher's or author's thought. Certainly there is a case to be made for this with regard to the process of deterritorialization, where something is wrested from a territory and deterritorialized upon a new territory, like the animal paw that is deterritorialized from the earth and reterritorialize on the branch. however, it seems to me that the more interesting aspect of Deleuze's approach to other philosophers and art is not so much his "monstrous creations" (i seldom find them particularly monstrous), but rather their Gaussian or Riemannian style, where he explores them in terms of their own internal organization and metric, without reducing them to something alien such as history, society, biology (reductivism), or the signifier. What we find in Deleuze's approach to phenomena is a Gaussian technique. for instance, take Deleuze's books on cinema, his book on Francis bacon, or his study of sacher-masoch. In the first instance, Deleuze's carefully separates cinema from narrative and the signifier, studying it in terms of its specific organization pertaining to the production of images. in the case of bacon, Deleuze doesn't look for an underlying narrative or "meaning", but instead studies the manner in which bacon composes and

organizes his images and lines, both in terms of their production and actuality, so as to liberate a logic of sensation. finally, in approaching sacher-masoch, Deleuze allows Sacher-Masoch's novels to speak for themselves in terms of their desire and relation to pain, stalwartly refusing to reduce masochism to the complement of sadism. in each case, we have a local exploration of a "space" of multiplicities that is extremely precise. Sean Bowden the priority of events: **Deleuze's logic of sense sean Bowden, the priority of events: Deleuze's logic of sense, Edinburgh university press, 2011, 296pp., $40.00 (hbk), isbn 9780748643646 Reviewed by David Scott, Coppin state university and Gaussian spaces and multiplicities a note on Gaussian spaces and multiplicities posted by larval subjects under uncategorized.** What you "see" is not always what you "get." Many people mistakenly take their own visions literally without "seeing through" the various possibilities that are outside their belief system, knowledge or skill base. Deeper reality is not remote in the physical sense but in a psychological sense. The archetypes of the collective unconscious are arrayed behind the scenes of current worldwide conditions, of crisis and confusion. They mirror our own states back at us, whether we perceive them as such or interpret them plausibly or not. The noise of ordinary consciousness and beliefs drowns out the signal. Unconsciousness is the background of our ordinary awareness. Our organism is very much at the center of such effects. The organismic source is our human bodies and the focus of human consciousness. The fantasy principle dethrones (e) reality, but can be dissociative or compensatory. The human mind is (=) a meme-Scape. Pre-conceived concepts vie with (e&eb) structures, concepts with images. Like scientists who ignore (e) assumed truths, we leapfrog over (e&eb) our beliefs and personality deficits, claiming (eb) idiosyncratic imagination is (=) literal reality. It couldn't be further from the truth and symbolism is utterly lost. The metaphor that might heal us enslaves (e) us. Perhaps images like the holographic universe have an implicate order. Can we have a sense of the cosmos in the world without projecting myriad fantasies on it that we embrace literally? Has the world become so horrible it is unreasonable to be realistic? We may need to look at our drives and wishes, rather than the fantasy content. Psyche constructs reality. Our

experience of so-called reality is always mediated by our image of it. Even if all the contents of the psyche are real, that doesn't mean they are realistic. That psyche is real is still a radical proposition, but psychic politics certainly color the self-image and ideas of everyone. We observe and participate with images. It is not a question of nature or nurture (genes alone or experience alone). Rather, everything is (=) both. We inherit (e&eb) the structures that make our experience (=) what it is. But the structure itself is "empty," and each human culture "fills" it with its own specific adaptations. It is difficult to define an archetype and set boundaries that distinguish it from others. In a hologram each part contains all the information but in lower resolution. Archetypes have this holographic quality. There are patterns within patterns within patterns. Some overlap with others, and some are nested inside others. Archetypal realities, passed on through DNA, are expressed in distinctive neuronal tracts in the brain. They include customs and laws regarding property, incest, marriage, kinship, and social status or roles; myths and legends; beliefs about the supernatural; gambling, adultery, homicide, schizophrenia, and the therapies to deal with them. A mythic and visionary language of immediate experience encompasses themes of deepest, highest, and ultimate concern. Most fantasy-based individuals are at a complete loss to coherently explain their own conventional behavior much less anomalous events and their deep meaning, much less the cultural unconscious or mythological unconscious matrix. But they try, and become utterly entrenched in their belief that they are right about the nature of the world and reality. We have pseudo-memories about (e&eb) our personal lives. Why not more so for our collective life? The subject matter often revolves around catastrophe, creation and the mythopoeic forces of mankind. Ignorant of such dynamics, interpretive mistakes and displaced psychic contents proliferate into errors of fact. Propaganda, media distortions, memes, and disinformation compound the social problem of misapprehension further. Shameless self-promotion by personalities of such ideas leads to cults. They make up myths about the myths of by-gone eras. Roiling unconscious images can be fatally confusing. Thought illusions culminate in projections and projections of mythology. Jung suggested symbols live only as long as they are pregnant with meaning. Philosophy arose from

criticism of myth, from discussing and challenging it. In science, we criticize, reject and eliminate theories. At the edge of the abyss of the unknown, new signs and symbols emerge. Credible theories and paradigms must include biology, physics, and neurophysiology. One of the reasons people "see God", or a guru, or anomalies may be because our brains are constructed to see reality through the eyes of others. There are heaps of mirror neurons which are there to make us feel the 'other'. Mirror neurons do for psychology what DNA did for biology. They provide a unifying framework and help explain a host of mental abilities. As in the psychochemical processes of empathy or falling in love, a complex feedback loop sustains a state of mind. But when we empathically transpose ourselves into someone else's position, we expose ourselves to that reality -- cognitively and emotionally. The unconscious complicates empathy, both ways. Mirror neurons might well play a role in bonding, language and self-awareness. Naively, we take too much as self-evident. But 'seeing' does not always 'believe', though many make this error or leap in logic and formulate their choices and future accordingly. Yet, there is only one way to learn what consciousness is. Experience. But we have no satisfactory explanatory edifice for consciousness. Would such a theory release in each of us our own inner knowledge of the creativity of our own consciousness, and its infinite possibilities? The problem is trying to define a verb, a dynamic, as if it were a noun. But we do recognize the effect of consciousness. It functions to mediate states of consciousness, high and low psychobiological arousal. Consciousness is the subconscious lifted up by the physical body. When the body fails, the consciousness collapses back into the subconscious. All our thoughts come from the subconscious which can see our intentions but not our world. This relates somehow to intention being imaginary and not of the physical frictionized world (King). Gerald Edelman postulates that the flows of information in the brain are mediated through're-entrant' feedback loops. As evolution provides new cognitive functions, new re-entrant loops are established. Even language itself is an archetype -- a chaotic field of dynamic associations. A subtle net of tropes, grammar, symbols, and meaning, the program language begins in limbic resonance. Some phenomena generate their own language patterns,

nomenclature, and internal coherence of meaning and representation. n a holographic universe, even time and space could no longer be viewed as fundamentals. Because concepts such as location break down in a universe in which nothing is truly separate from anything else, time and three-dimensional space would also have to be viewed as projections of this deeper order. At its deeper level reality is a sort of super hologram in which the past, present, and future all exist simultaneously. This suggests that given the proper tools it might even be possible to someday reach into the super holographic level of reality and pluck out scenes from the long-forgotten past. Or not. A fantasy of such penetration or phenomenon inside the head is not the same as that penetration. **Jung in the 21st Century: Synchronicity and science By John Ryan Haule Mind Control Countermeasures.** Adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrograde, inimically inverse, violatively unsimilar, diversely dissimilar some statements in Vedanta, it is not superfluous or redundant. It is evolutionary. Authetication determination, unimpeachable validation, incontrovertible establishment, apodictic evidence, reliable roll out, unfailing cinquecento **quattrocento trecento,** incontrovertible indication of the evolutionary process towards understanding is ever present in all, its thematic and discursive form. it is the relentless holding on to a theory this is despicable albeit the theory is not. Look at this paragraph. It is the 'I am the body/mind' belief that gives rise to the 'I am not (e) the world' belief. These two beliefs are (=) co-created."Consciousness projects (e&eb) the appearance of the mind, body and world by taking the shape of thinking, sensing and perceiving." "Attention is Consciousness with an object. When the object vanishes, attention simply remains what it always is, Consciousness.""There is (=) no purpose to (e) meditation. The purpose is already accomplished." "Everything that is experienced is experienced by, through, in and as Consciousness." "The seen cannot be separated from seeing and seeing cannot be separated from (e) Consciousness.""The Reality of any experience is not hidden in the appearance; it is expressed by the appearance." "Once we see that everything is Consciousness... Maya still dances, but it is a dance of love not seduction." There are some truly excellent sections

such as a long one taking, as a starting point, the following quotation from the painter Cézanne: "Everything vanishes, falls apart, doesn't it? Nature is always the same but nothing in her that appears to us lasts. Our art must render the thrill of her performance, along with her elements, the appearance of all her changes. It must give us a taste of her Eternity." Rupert begins: "That statement must be one of the clearest and most profound expressions of the nature and purpose of art in our era." And he goes on for some 14 pages to elaborate on this claim, examining the nature of the 'elements' Existence, appearance and Consciousness and their relationship

Some of the following concepts are modeled; some are not. For those which are not modelled kindly see one of the papers in the series. This note has the axiomatic predication to give a holistic view of the study. Here we talk of the characterstics of systems which satisfy the condition of cosmological constant or any other for that matter or variable or axiomatic predication. . There are lots of zeroes corresponding and concomitant to the second law of black holes. Infact as many as that of extant and existential black holes exist. At the outset, it is to be stated that there are hadrons in every system. Supersymmetry between forces and matter, with both open and closed strings; no tachyon; group symmetry is SO (32) and its axiomatic predications, predicational anteriorities, character constitution shall be extant and existential in very many systems, and the characterstics are taken in to consideration in the classification scheme. Many systems have such fundamental instabilities like that of quantum gravity and characterstics of those systems form the citadel and fulcrum, bulwark and manor, mainstay and reinforcement, alcazar and chateau theory has a fundamental instability on which the classification schémas are valid. There are lots of systems which follow the axioms of string theory. It is the characterstics of this system which are taken in to consideration in the classification scheme. "All institutionalization involves common moral as well as other values. Collectivity obligations are, therefore, an aspect of every institutionalized role. But in certain contexts of orientation-choice, these obligations may be latent ...." (Parsons, 1951, p. 99). There are various systems that have the same bastion, support system, stylobate and sentinel as that of the Deleuzean terms and predications and phenomenological methodologies systemized. Look at this beautiful passage whose stability analysis would be eye opener. To "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new concepts and meanings that could not be included in the old system. At the heart of Western metaphysics, for example, Derrida finds the opposition between "speech" and "writing." This binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior In Derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), Derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. **THE POSTMODERN TURN IN PHILOSOPHY: THEORETICAL PROVOCATIONS AND NORMATIVE DEFICITS by Steven Best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html** Each and every system has electrons, neutrons and protons and for that matter quarks. There shall be strong nuclear force and weak nuclear force. There are many systems that satisfy the criterion specified by the equation, principle or statement in question. Characterstics of the investigating systems form the bastion for the classification scheme and doxa thereof. Systemic differentiation is conducted. Despite gravity being constant, there exists gravity between two objects, and this could be taken as a system. Depending upon some parametric representationalities, functionalities, advantageousness, appropriateness, benefit,

facilitation, fittingness, helpfulness, instrumentality, merit, practicality, serviceability, suitability, and usefulness, utility, these systems could be classified in to various categories. In respect of an equation, there shall be many systems that satisfy the given equation. Equations themselves could be by the utilization of the model solved term by term as has been exemplified and illustrated many time in the previous papers. There is lot of systems that could be brought in to the orbit of and gamut of the theory in question which the investigatable systems satisfy the axiomatic predications and postulation alcovishness of the systems in question. Towards the end of classificational consummation, consolidation, corporation and concatenation we take the characterstics of the systems, the predicational interiorities, ontological consonance and primordial exactitude, accolytish representations, functional topology, apocryphal aneurism and atrophied asseveration, event at contracted points, and other parameters as the bastion and stylobate of the stratification purposes. Such totalistic entities would have easy paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of (e) homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, projection and prognostication as consideration of the investigatory systems. Any scale can be used that is convenient to the classification scheme. It is important to note that the scheme of classification and the stratification doxa must not be adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatiously gainsaying or conformingly pugnacious to the axiomatic predications and postulation alcovishness of the theory in question or the equation representative or constitutive thereof. Sole intention, main objective and primary aim is twofold. One is towards the end of circumvention of the extra equation and the concomitant and corresponding variable therein. Second is avoidance of clustered congest, swarmed huddle, mustered pack, sardine squash, and swamp throng in the scheme of classification. Only thing that is sought out is the consonance in the entire diaspora and body fabric of the systems under study. This statement is true and holds unmistakably true for all the papers and I sincerely entreat readers to remember the statement and read the paper against this background. When we write A+B we mean by that B is being added to A or vice versa. It is like adding milk to water and water to milk. There may or may not be a time gap. As said earlier there may be many systems that satisfy the conditions of the equations and those systems that are investigatory or investigatable are taken in to consideration based on their characterstics in the classification doxa. When there are more than two entities, we can take logarithm and find the value of that factor to be taken with anti log to obtain prediction and projected values of the model. In case of A-B, we are removing B from A and that means B is eating up A. These factors are taken in to consideration in the application of the model to equations. Log (ab) and Log (a+b) is well defined. Model stands out as universal testament and template for application to each sentence and equation what with the quantification process done and the correlations well defined. In the eventuality of non existence of any connection at some phase, the model would have the accentuation and attrition coefficients and detritions coefficients as zero rendering the equations of concatenation simpler. Projection formula which incorporates in its diaspora the initial values provide autheticative determination, unimpeachable validation, incontrovertible establishment , apodictic evidence, reliable roll out, unfailing cinquecento quattrocento trecento, incontrovertible indication of the final finale, notwithstanding appellation, appellative, brand, cognomen, compellation, designation, flag, handle, identification, label, moniker, nomen, slot, style, surname, tab, tag, term, title designation, appellation, appellative, class, classification, denomination, description, epithet, of the classification scepter, scenario, scimitar, schottische. There is pure and impure consciousness in every one. Gratification producing and deprivation producing one's can be easily classified

from individual general ledgers. Similar analogy holds for collective general ledger and cosmic general ledger or nature's general ledger. It is also to be noted that while dealing with equations towards the end of consummation of the measure, it is necessary that the two variables are to be classified in to three sections and each one would have the adventitious and decidedly stated relationship whereby the fundamental equations are drawn up and the analysis made. All the parametric representationalities, conditionalities, orientationalities remain unequivocal as stated in the variables stated in to consideration section and are different from module to module. In essence the paper is to be read as holistic one with the sole intention, primary objective and raison d' etre to build a TOE. Towards the end of circumvention of typing hundreds of superscripts and subscripts which would be a sardine squash and the concomitant operational difficulties, model is presented in piece meal of nine modules. Logarithms are to be taken in respect of those which incorporate more than one variable in bra-ket. Values of log (ab) and log (a+b) are readily available. Anti logarithm shall be taken at the end to predict, project, prognosticate the value of the variable. This is true for tensors, vectors and other variables too. Affirmational assertion, and explanation justification, statement of vindication, annotational commentary, and explicational glossary for each and every system changes and physical interpretation of results is one thing that is to be with earnest endeavour and feverish and febrile expediency. Editions never ever mean the same and identicality of parameters and this has been explicitly and unmistakably stated in the model apriori itself. Akin and analogous, cognate and concurrent, correspondent and congruent, comparable and complementary, synonymous and duple, tantamount and agnate, commensurate and correlative representation is only to highlight the importance and subterfuge the replication of work. It is my fervent solicitation to kindly bear with me for any lapses, notwithstanding the orchestrated efforts for a paper without any minor errors. Postulation predication, conclusive presumption, differential presuppositions, underscored decidedly axiomatic statement of the statement, equation form the bye word or the watch word in the aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement-argument, articulation, assertion, asseveration, choice of words, commentary and communication, declaration and definition, delivery-diction, elucidation-emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation tone and tenor of stratification. Now, how do find the reaction of systems to these singularities. You do the same thing a boss does for you. "Problematize" the events and see how you behave. I will resort to "pressure tactics". "intimidation of deriding report", or "cut in the increment" to make you undergo trials, travails and tribulations. I am happy to see if you improve your work; but may or may not be sad if you succumb to it and hang yourself! We do the same thing with systems. systems show conducive response, felicitous reciprocation or behave erratically with inner roil, eponymous radicality without and with blitzy conviction say like a solipsist nature of bellicose and blustering particles, or for that matter coruscation, trepidiational motion in fluid flows, or seemingly perfidious incendiaries in gormandizing fellow elementary particles, abnormal ebullitions, surcharges calumniations and unwarranted(you think so but the system idoes not!) unrighteous fulminations. Amount of emotion or affect that is appropriate or expected in an given form of interaction. Particular individuals and diffuse obligations (see c and d) are associated with affectivity, whereas contacts with many individuals (universalistic) in a bureaucracy may be devoid of emotion and characterized by affective neutrality. Affective neutrality may refer to self discipline and the deferment of gratification (eg Weber's spirit of capitalism). In contrast, affectivity may be associated with expressing emotions. Adams and Sydie also refer to affective neutrality being associated with ego control (p. 15). (Parsons: Wikipedia). So the point that is made here is "like we problematize the "events" to understand the human behaviour we have to "problematize" the events of systems to understand their behaviour. This statement is made in connection to the fact that there shall be creation or destruction of particles or complete obliteration of the system (blackhole evaporation) or obfuscation of results. Some systems are like "inside traders" they will not put signature atoll! How do you find they did it! Anyway, there are possibilities of a CIA finding out as they recently did! So we can do the same thing witn systems to. This is accentuation, corroboration, fortification, .fomentatory note to explain the various coefficients we have used in the model as also the dissipations called door. In the bank example we have clarified that various systems are individually conservative, and their conservativeness extends holisticallytoo.that one law is universal does not mean

there is complete adjudication of nonexistence of totality or global or holistic figure. Total always exists and "individual" systems always exists, if we donot bring Kant in to picture! For the time being let us not! Equations would become more ensorcelled and frenzied..... philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". In consciousness research, two rival sets of theories can be recognized: (A) Scientific material interpretations of consciousness are based on axioms that view consciousness in the context of highly advanced intentional processing of information in which subject-object relations evolve, and (B) humanistic interpretations of consciousness are based on axioms that view consciousness in the context of, say, "centered pulsations" that enable a conscious agent to act from his or her center of awareness. In this paper I will argue for the selection of axioms that favor humanistic interpretations of consciousness. The Pursuit of Autonomy Interdisciplinary Observations to Human Consciousness: Wautischer, Helmut (1993) the Pursuit of Autonomy. Interdisciplinary Observations to Human Consciousness: Social Neuroscience Bulletin, 6 (4). pp. 52-56.Everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality; **CONSCIOUSNESS AND ITS UNIQUENESS June 12, 2012 at 8:40pm (See Deleuze Logic of sense, Wikipedia and Stanford encyclopedia for more details)** Atrocious contrivance and device, gratificational primogeniture, calamitous dodge and expedient, depraved, destructive, disastrous, execrable gambit and gimmick, great solace and succor, iniquitous, injurious, loathsome, low, maleficent, malevolent machination and maneuver, promethaleon of candor, frankness, honesty, honor, ingenuousness, innocence, openness, reality, sincerity, truthfulness, spiteful, stinking, ugly, unpleasant, unpropitious play and ploy, progenitor of jurisprudence and circumspection, racket and ruse, savvy, scam, stratagem, subterfuge, tactic and wile, forthrightness, honesty, truthfulness embodied and personified is the structural predisposition and dispensation of the complex. Nobody need have to take the Brahman-Anti Brahman agency to be a recalcitrant, repugnant, and refractory proposition. God gives and also takes away. That is the point made. ‘

And of course lesser gods, and God doth follow them and they follow God. That is why they are lesser gods. And wield the power on lesser mortals like us. When we mean gratification and deprivation we mean celestial Brahman Anti Brahman and Terrestial Brahman and Anti Brahman agencies producing such happiness or sadness by any means such as (blending- amalgamation, adulterant-adulteration, amalgam-amalgamation, blend- combination,

composite- compound, debasement-denaturant, fusion-hybrid, intermixture-reduction, amalgam- mixture, composite- compound, fusion-mishmash, aggregate- alloy, amalgamation-blend, commixture-composite, composition-compost, conglomerate-fusion, goulash-medley, mishmash-stew, blend-coalition, commixture-compound, federation-heating, integration-junction, melting-merger, synthesis-unification, assortment- combination, adulteration-alloy, assimilation-association, batter-blend, brewed combine, composite compound, concoction confection, conglomeration cross of money, orientation(dis), penury creation, murder, mayhem plunder, pillage, apocalypse, Armageddon in various forms and permutation and combination of Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness: with each action accounted and showing balance in the nature's general ledger as also actionlessness:e). All these are recorded in Nature's general ledger. Neuron DNA encapsulates all the actions passions, actions, interactions, transactions or the lack of it at various levels of individual, collective and cosmic levels. In celestial Brahman Anti Brahman ledger there is cessation of dualities. One more point to be made is if 5=3+2, then we can say 5 gormandizes 3 or 5 gobbles up 2. Quintessentially such a statement arises out of the difference of LHS and RHS term wise. It must be remembered despite all theoretical abstractions and generalizations be it proton mass or Higgs Boson, or topology of a space, it is notwithstanding the fundamentality of nature is just a number. Just everything in Physics and mathematics is just a number. Let us assume I buy rice in Bangalore (District) worth 5k.g. I can always divide 5k.g.s of rice in categories of Bangalore

(South), Bangalore (west), Bangalore (East) and Bangalore (South). On similar basis, there are billions of neutrons, protons and concomitant interactions which would satisfy the given numerical value. Example of gravity above makes things clear. Everything atleast is a function of time and someplace or some other coordinates, subordinates or superordinates. Mass of proton is 2.47 does not mean total mass of all protons is 2.47. It is this argument that is used in the classification scheme. By the very nature, self the unmonitoring but the witness consciousness agency identifies itself with something, be it a profession, an actor on the screen, when it does not have a identity. Such examples are legion, a Charlie Chaplin in "The Kid" or Raj kapoor in "Jagte Raho". These are people who just wanted to live and living without being put under surveillance every moment is the greatest part of their life. This is "Absolute Subjectivity" of Shiva. Such a state is always craving for identification for in pure consciousness all identities are lost. It is this "identitylessness" or having very minimal identity is what pervades most in India and this is due to expansion of individual consciousness. When that happens, be it education, or wealth, the "identity' becomes so well entrenched, they become the exact opposite of "pure consciousness"," impure consciousness" (Shakti). This is a state of "Absolute or relativistic objectivity". There was this case of a lady who went to such an extent of planning, she was sure that whatever she does must work out. When it did not, she lapsed in to "identity crisis". Such thing as this might happen in well developed countries where the competition is extremal and to be achieved at any cost. Realisation that one is not the doer but the witness consciousness is just like watching a movie in which one has acted. It is in this sense "advaita", "non duality' seems to be of cardinal and preeminent importance for the analytical treat ment of "identity' and "identity lessness", both happen in space time. "Multiverses" in this sense is also like a different film of an actor released at different theaters and it is in this sense we state that that every object in space time is a multiverse. You must have totally immersed in doing a "string theory" problem or watching a "film", oblivious of surroundings. Essentially creation (space time) also is what is happening in mind and honed in to you billion times. Identification of anything with a designation or signification is one that leads to

association within spacetime, which would only have pathological consequences and detrimental ramification as stated above or positive implications of "happiness" because of the identity itself. Search or 'show" has to go on! It gives a sense of consummation. a watching pot never boils; donot go to check nature; you are disturbing it; it will never show its true colours bhakti and Virakthi, gratification and deprivation, raga and dwesha is essential to live in space. It is the quantum interest thereof that produces energy to live .gratification and deprivation which form the bastion, pillar, post and stylobate and sentinel of all human actions depends on tams, rajas, and sattva of individual or insentient object. Verily tamas, rajas and satva are the characterstics of space itself, nay the behavioural pattern and attitudinal orientation. There are peaceful places like Vienna or Sringeri and there are violent places like others. Space is warped and woofed by the Akshara, the quantum information of primordial syllable Om. cosmic general ledger is the one vibration of which is brahman; quantum information possessions thereof form the sentinel of all transactions (Gargi to Yajnavalkya: my interpretation)adhocism, nonseqientiality of events, anti sequentiality of events, simultaneity of phenomenon, concurrent happenings , improbability of phenomenon happening itself are characterstics of brahman-anti brahman (Terrestial and celestial) are the cause of non predictability of happenings like in black hole, or in the areas of operation of brahman-anti brahman (Terrestial and celestial). brahman AntiBrahman actions, interactions, decisions, circumspection, adhocism, satisfies sensitive to initial conditions; topologically mixing; dense periodic orbits, the axiomatic predications and postulation alcovishness of chaos theory. Like gravity is a characteristic of space, so is tamas, rajas, sattva, with spaces knotted together by quantum information (for models see below) .accidents are preplanned actions of Terrestial and celestial Brahman AntiBrahman. All actions are perpetrated by them, be it train on fire or death. Belief systems play an important role in the determination of tamas, rajas, and satva state of human beings. There is neither cause nor effect. All actions are performed by Terrestial and celestial Brahman anti Brahman and it is off the record. psychic energy is negative equivalent to –mc^2, so that second law of thermodynamics is satisfied on a holistic basis vis-à-vis material energy

ParaBrahaman is infinite individual, collective and cosmic consciousnesses are infinite (concomitant general ledgers) for they are in the evolution mode with respect to each other in that order. Later emanates from the former as transactions emanate from the conservative state of assets and liabilities. Is not zero true for infinite conservative systems inside space and time? zero is the only consummation, consolidation and conservation. if fifty out of seventy members come out of the room at whatever time you come out, talk about the grand limbo of oblivion and hibernation of lifeless world you are going to and if you are not a celebrity, then it is not coincidence, it is coordination. Nature is playing the same trick on you. That a generalised law is there and you can measure does not make things less good albeit recalcitrant, refractory and rogue particles occur time and again, by and large. Entering in to all living beings through his constant part on the new moon night, he is born therefrom next morning. Absolute ParaBrahaman (zero) projects himself again and again on the screen of individual consciousness which appears of depairing despondent world. Some where the balancing act is done on karma. cosmic general ledger reflects upon itself to keep all transactional ties conservative (Brihadarnkayaka Upanishad: 1.5 14) for models see next paper through its divine power ( i interpret man creates god; it is attribution and ascription of various Shaktis would the formation of god takes place like space and time) the self assumes (eb) the form of a deity man can contemplate and venerate, even though siva, the pure subject, can never in fact be an object of meditation (adhyeya). Until we realise our true identity with Sankara, he is worshipped and conceived to be a reality alien to ourselves. While we are in the realm of creation, he too is a creation or mode or appearing of the absolute, manifest to us in meditation, through his freedom as an eternal, omniscient being. There is no gulf between the created and the uncreated creator: nothing in reality, although an object of knowledge, ceases to be Siva: this is the reason why meditation [on this or that aspect] of reality bestows its fruit. The world of the senses and mind appears to the well awakened (suprabuddha) as a theophany, an eternal revealing of god in his creation. The doctrine of vibration declares that "there is no state in word, meaning or thought, either at the beginning, middle or end, which is not Siva." 59 to utter any word is, in

reality, to intone a sacred formula.every act is a part of Siva's eternal cosmic liturgy, every movement of the body a ritual gesture (mudras), and every thought, god's thought. By what path are you not attainable? What words do not speak of you? in which meditation are you not an object of contemplation? What indeed are you not, o lord? 61 spandashakti, which accounts for the appearing of all things, is also the means by which deity in its many varied forms appears to man. Ksemaraja concludes: the ultimate object of worship of any theistic school differs not from the Spanda principle. The diversity of meditation is due solely to the absolute freedom of Spanda. 62 Sankara is not only the supreme object of devotion; as the static polarity of the absolute, he is the inner reality which holds together its Siva and Sakti outer manifestations. Phenomena are patterns of cognitions projected onto the surface of self-luminous siva-consciousness. There they become apparent, directly revealed to consciousness according to their manifest form. Siva is accordingly symbolized as the ground or surface of awareness, smooth and even like a screen (samabhittitalopama). 64 inscribed on this screen (kutfya) are the countless manifest forms which appear within it rendering it as diverse and beautiful as a fossil ammonite (falagrama). 65 Siva is the sacred ground upon which the cosmic mandalas are drawn, the absolute surface of inscription which bears the mark (cihna) of the universe. Abhinava writes: the variety of this world can only be manifest if the highest lord, who is essentially the pure light of consciousness, exists; just as a surface is necessary for a picture. if external objects were perceived in isolation then, because 'blue' and 'yellow', etc., are self-confined and the perceptions [we have of them] refer to their objects alone and so are insentient, mute and dumb in relation to one another . . . how would it be possible to be aware that an object is variegated? but just as depths and elevations can be represented by lines on a smooth wall, and we perceive [a female figure and think], similarly it is possible to be aware of differences in the variegated (contents of experience) only if all the diverse perceptions are connected together on the one wall of the universal light of consciousness. 66 Siva is the perfect artist who, without need of canvas or brush, paints the world pictures. The instant he imagines it, it appears spontaneously, perfect in every respect. The colours he uses are the varying shades

and gradations of his own Spanda energy and the medium his own consciousness. Intractable disaffect, Mutinous obdurate, Rebellious recalcitrant, Seditious refractory, Insubordinate insurgent, antagonistically antipodean, discrepant incompatible, Contumacious coadjutant, Conflicting converse, Unyielding obdurate, the antagonistic and anachronistic dispositional ties of Anti Brahman are in fact the sine qua non of Brahman only (See Vishwaroopa of Mahavishnu, which contains the polarised positives on RHS and negatives on LHS). The universe is coloured with the dye of its own nature (svabhava) by the power of Siva's consciousness (a//). Rajanaka rama says: homage to him who paints the picture of the three worlds, thereby displaying in full evidence his amazing genius {pratibha)\ to Sambhu who is beautiful with the hundreds of appearances laid out by the brush of his own unique, subtle and pure energy. 68 analogously, at the microcosmic level, all the cognitions and emotions, etc., which make up the individual personality form the outward flow of essentially introverted consciousness. They are specific pulsations (vise$aspanda) or aspects of the universal pulsation (samanyaspanda) of pure t consciousness. at the lower level, within the domain of maya, they represent the play in the fettered soul of the three primary qualities (gunas) or feeling-tones' which permeate to varying degrees his daily experience. These are: 1) sattva — the quality of goodness and luminosity which accompanies blissful experience both aesthetic and spiritual. 2) Rajas — the passion or agitation which oscillates between the extremes of 'light' and 'darkness' and characterises inherently painful experiences. 3) Tamas — the torpor and delusion which accompany states of inertia and ignorance. 69 the liberated soul recognises that these three are the natural and uncreated powers of pure consciousness. for him they are manifest respectively as: 1) Sankara 's power of knowledge (jnana) — the light of consciousness (prakasa); 2) the power of action (kriya) — the reflective awareness of consciousness (vimarsa)\ 3) the power of maya — which does not mean here the world of diversity, but the initial subtle distinction which appears between subject and object in pure consciousness. 70(doctrine of vibration: mark S.G. Dyczkowski) .to "deconstruct" is not the same as to destroy. Deconstruction attempts to undo logical contradictions, to overturn rigid conceptual oppositions while releasing new

concepts and meanings that could not be included in the old system. at the heart of western metaphysics, for example, derrida finds the opposition between "speech" and "writing." this binary logic functions in an illicit way to establish speech as the means of giving "presence" to the world, while writing is deemed derivative and inferior in derrida's sense of "grammatology," however, all production of meaning is writing and subject to the infinite play of signification. By taking away the transcendental signified and advancing the concept of "differance" (language organized around difference and deferred, or mediated, understandings), derrida, like Nietzsche, wants to leave us without transcendental illusions, metaphysical unities, and foundations that constrain thought and creativity. The postmodern turn in philosophy: theoretical provocations and normative deficits by Steven best and Douglas Kelner http://www.gseis.ucla.edu/faculty/kellner/kellner.html. You cannot see beyond consciousness for that is the limit of your storage of information or the basis for understanding further information. Turiya is empty bucket; it is from which prajna the measure of  arises; turiya measures the number of quantum information stores in other three states of waking, dreams, and dreamless deep sleep (Manduka Upanishad; my interpretation for models next paper).let there be path where angels tread: it is a force never born never dies the brahman an anti brahman verily there are parallel universes in (eb) in the ambit of space and time; the society of sinners and the sinned with the grammar of evil; parallel universes contain(e) multiverses in them; all physics, philosophy,  break down(e) in these universes; astronomical gratification ,deprivation, glorification, mortification, projections, axes, dimensions, rotations, foldings, concurrences, motions, shocks, are the axiomatic predications of multiverses and parallel universes; death is when self (the witness consciousness; register of individual general ledger) leaves the body; ostensibly parallel seem to have (e) an objective; parallel universes donot (e) have objectives; their paradigm of relational content, differentiated system of expressly oriented actions with primary focus and locus of homologues receptiveness and differentially instrumental activity, variable universalism and particularism, imperative compatibilities and structural variabilities, interactional dynamical orientation, institutionalization and internalisation of pattern variables common

attitudinal orientation of constituionalisation of internalized dispositions, and a qualitative gradient of structural differentiation and ascribed particularistic solidarity abstraction or interactional dynamics, internal differentiation, structural morphology, formal characterization, concept formulation, phenomenological methodologies, constituent structure, transformational minimal condition, paradigmatic feasibilities, programmatic plausibilities, comparative variability, normative aspect of expectational prediction, affection- appreciation, awareness- discernment, emotion-feeling, gut reaction, heart insight, intuition, judgment, keenness, perceptiveness, rationale, sensation, sense, sensitiveness, sensitivity, sentiment, susceptibility, taste, vibes, admiration, aesthetic sense, affection, appraisal, assessment, attraction, awareness, cognizance, commendation, comprehension, enjoyment, esteem, estimation, grasp, high regard, knowledge, liking, love, perceptional realization, recognitional regard, relish, respect, responsiveness, sensibility, sensitiveness, sensitivity, sympathy, understanding, valuation, acknowledgment, airing, baring, , confessional defenselessness, denudation- denunciation, disclosure, display, divulgence-exhibition, hazard, introduction, jeopardy, laying open, liability, manifestation, nakedness, openness, peril, presentation, publicity, revelation, risk, showing, susceptibility , susceptiveness,susceptivity, unfolding, unmasking, unveiling, vulnerability, vulnerableness, action, affection, appreciation and ardor, behavior, capacity, compassion, concern, cultivation and culture, delicacy, discernment and discrimination, emotion, empathy, faculty, fervor, fondness, heat, imagination, impression and intelligence, intensity and intuition, judgment, keenness- palpability, passion, pathos, pity, reaction, refinement, sensibility, sensitivity, sentiment, sentimentality , sharpness, spirit, sympathy, tangibility, taste, tenderness, understanding, warmth, fond memories, hearts and flowers, homesickness, longing, pining, reminiscence, ,schmaltz, sentimentality wistfulness, yearning satiation is only ostensible; their job is to get the job done; their business is to mind other's business ;all laws of physics and finance breakdown; it is antilaw,piscatorial Piratish and bubonic bucaneerishness; there is free transactions between in multiverses; when reflected on the screen of consciousness they appear as extant existential universe. Transactions are astronomical and gratification

deprivation, glorification and mortification incalculable; they have no barriers of region, religion, language which they create for others. theirs objectives are: corruption criminality, debasement debauchery, degradation evil immoral lewdness licentiousness, perversion sinfulness, vice viciousness, vileness wickedness, improbity corruption, atrocity decadence, degeneration degradation, depravity evil, immorality impurity, infamy iniquity, looseness lubricity, perversion profligacy, sinfulness turpitude,vice viciousness, vulgarity wickedness,disgrace bad reputation, abasement abuse, baseness black eye, blemish blur, brand comedown,contempt contumely, corruption culpability, debasement debasing, defamation degradation, derision disbarment, discredit disesteem, disfavor dishonor, disrepute disrespect, humbling humiliationignominy ill repute, infamy ingloriousness, meanness obloquy, odium opprobrium, pollution ,put-down reproach, scandal scorn, slander slight, slur spot, stain stigma, taint tarnish, turpitude venality, enormity Horribleness , abomination atrociousness, atrocity crime and the positive aspectionalities thereof .let there not be any incursions where evil mercenaries in the garb of mendicants thrive with adversely antagonistic, inexorably irreconcilable, antithetically antipodal, diametrically opposed, repugnantly retrogressive, inimically inverse, violatively unsimilar, diversely dissimilar, antipathetically anti thetical, conflictingly combative, obstructively pugnacious, inimically obstructive, repellently restrictive, disputatious gainsaying or conformingly pugnacious Manichaean artifice, agathokakological malevolence, agley chicane fourberie, fraud, furtiveness, gambit, awry and bad ,flagitious mechanizations, cacodemonic gambit, hanky-panky, intrigue, machination, deprecatory and diabolic underhandedness and wiles, energumenical skullduggery, execrable and goetic chicanery, lenocinant stratagem, malefic and maleficent strategy, malominous sell, sellout, sham subterfuge, peccable, humbug, imposture, misrepresentation peccant fraud, fraudulence, hoax, qued , scelerate pretense, ruse, sell, sellout, sham stratagems' of maladroit manouevures , sinistral ,tortious and venal sophistry, contingency enterprise, endangerment emprise, jeopardizing adventure, jejune jeremiad, avoidable inertia, latent passiveness . statement for

Mephistophelean mercenaries philosophy merges with ontology; ontology merges with univocity of being; analogy has always a theological vision; not a philosophical vision; one becomes adapted to the forms of god; self and world; the univocity of being does not mean that there is one and the same being; on the contrary, beings are multiple and different they are always produced by disjunctive synthesis; and they themselves are disintegrated and disjoint and divergent; membra disjuncta.the univocity of being signifies that that being is a voice that is said and it is said in one and the same "consciousness". everything about which consciousness is spoken about being is the same for everything for which it is said like gravity; it occurs therefore as an unique event for everything; for everything for which it happens; eventum tantum; it is the ultimate form for all of the forms; and all these forms are disjointed; it brings about resonance and ramification of its disjunction; the univocity of being merges with the positive use of the disjunctive synthesis, and this is the highest affirmation of its univocity like gravity; it is the eternal resurrection or a return itself, the affirmation of all chance in a single moment, the unique cast for all throws; a simple rejoinder for Einstein's god does not play dice; one being, one consciousness, for all forms and all times. a single instance for all that exists, a single phantom for all the living; a single voice for very hum of voices; or a single silence for all the silences; a single vacuum for all the vacuums; consciousness should not be said without occurring; if consciousness is one unique event in which all the events communicate with each other; univocity refers both to what occurs to what it is said. This is attributable to all states of bodies and states of affairs and the expressible of every proposition. So univocity of consciousness means the identity of the noematic attribute and that which is expressed linguistically and sense fully. Univocity means that it does not allow consciousness to be subsisting in a quasi state and but expresses in all pervading reality. Negative energy stabilising wormholes are considered candidates for (er) darkmatter. An analysis of negative energy might be a useful ingredient of the thoughts of travel at speeds more than the velocity of light. Peirce explains that an "imperfect continuum" possesses "topical singularities, or places of lower dimensionality" which break continuity, while "the parts of a perfect continuum have the same dimensionality

as the whole." (CP 4.642) Circular continua without singularities are time and space. (CP 6.210-212) By using math as his lens, Peirce envisioned singularities or lower dimensional holes in the imperfect cosmologic continuum, while maintaining that a probabilistic law of chance exists that entropy decreases **(Wikipedia).** On the side of philosophy**,** quite explicitly, Deleuze provides a critique of Husserl's "sleight-of-hand" change from a static categorical phenomenology to one supposedly more dynamically constitutive of reality. This a clear cut **chagrined circumvention**, **contraventional curbing**, **defeated disgruntlement**, **dissatisfactional downer**, **dragged failure**, **fizzled foil**, **grievance hindrance**, **impedimental irritation**, letdown, nonfulfillment, nonsuccess, obstruction, old one-two, **resent mental setback**, unfulfillment of the earlier theories makes Sartre contend that human existence is a conundrum whereby each of us exists, for as long as we live, within an overall condition of nothingness (nothing-ness)—that ultimately allows for free consciousness. But simultaneously, within our being (in the physical world), we are constrained to make continuous, conscious choices. It is this dichotomy that causes anguish, because choice (subjectivity) represents a limit on freedom within an otherwise unbridled range of thoughts. Subsequently, humans seek to flee our anguish through action-oriented constructs such as escapes, visualizations, or visions (such as dreams) designed to lead us toward some meaningful end, such as necessity, destiny, determinism (God), etc. Thus, in living our lives, we often become unconscious actors—Bourgeois, Feminist, Worker, Party Member, Frenchman, Canadian or American—each doing as we must to fulfill our chosen characters' destinies. However, Sartre contends our conscious choices (leading to often unconscious actions) run counter to our intellectual freedom. Yet we are bound to the conditioned and physical world—in which some form of action is always required. This leads to failed dreams of completion, as Sartre described them, because inevitably we are unable to bridge the void between the purity and spontaneity of thought and all-too constraining action; between the being and the nothingness that inherently coincide in our self. Deleuze insists on the absolute non-resemblance between what conditions (metastable, pre-individual plane of singularities) and what is conditioned

(actualized or the individuated 'thing'). Only then do "the conditions of the true genesis become apparent," Deleuze writes (Logic of Sense 105). Without addressing this issue, Deleuze's borrowings from Simondon, Lautman and structuralism remain untethered. Because of Bowden's decision to completely sidestep Deleuze's overt negotiations with phenomenology, his otherwise sophisticated commentary suggests a pedagogical narrowness. If, the omission of phenomenology strengthens Bowden's focus on the ontological priority of the event over substantial ontologies via the genealogy of figures and movements Deleuze outlines, it also incompletely articulates the problem of the sense-event as Deleuze formulates it. We have lost any idea of whom Deleuze takes to be his primary adversary, against whom he stakes out his position. This worry is most clearly manifested in Bowden's description of Deleuze's ontological prioritizing of events as offering a "transcendental ontology." For Bowden Deleuze's transcendental ontology defines events that "are ontologically prior to substances, 'all the way down'" (82). But Bowden provides a second definition: it is that the priority of events over worldly individuals means that the world of events "is not something external to the conditions of knowledge" (69). **SEAN BOWDEN the Priority of Events: Deleuze's Logic of Sense Sean Bowden, the Priority of Events: Deleuze's Logic of Sense, Edinburgh University Press, 2011, 296pp., $40.00 (hbk), ISBN 9780748643646. Reviewed by David Scott, Coppin State University Weltschmerz, agony, apprehension, blues, depression, dread, mid-life crisis, misgiving, nervousness, uneasiness, angst were expressed with such effulgent words…… aggrandizement-amplification, caricature and crock, understatement- unembellishment, elocution-emphasis, enunciation-inflection, announcement, argument, articulational assertion, asseverational choice of words, commentary and communication, declaration and definition, delivery, diction, elucidation, emphasis, enunciation-execution, explanation, exposition-formulation, idiom, interpretation and intonation about the human predicament is brought out in Being and Nothingness. Being and Nothingness offers a critique of Sigmund Freud's theory of the unconscious, based on the claim that consciousness is essentially self-conscious. Sartre also argues that**

**Freud's theory of repression is internally flawed. According to Sartre, in his clinical work, Freud encountered patients who seemed to embody a particular kind of paradox—they appeared to both know and not know the same thing. In response, Freud postulated the existence of the unconscious, which contains the "truth" of the traumas underlying the patients' behavior. This "truth" is actively repressed, which is made evident by the patients' resistance to its revelation during analysis. Yet what does the resisting if the patients are unaware of what they are repressing? Sartre finds the answer in what Freud calls the "censor". "The only level on which we can locate the refusal of the subject," Sartre writes, "is that of the censor.** The Holographic Principle, as conceived in current physics, applies to fields, and perhaps, even to more elementary entities, called strings and branes. These more elementary entities remain hypothetical at the time of this writing. There are many layers between the level of fields and that of neurons. We know that neurons are surrounded by a surface phospholipid membrane, which supports an electrochemical process that is fundamentally necessary to human experience. As the information supported by the membrane surface according to the Holographic Principle is proportional to its area, we should expect to find that the neural correlates of consciousness most clearly associated with those parts of the neuron that have the have a high proportion of neuronal surface area, and, in particular, receptive surface area. Dendrites account for an average of about 90% of the neuron is receptive surface area (Wong, 2002), and so would be expected to be the most important neural structures in information processes. We find this to be true **The Holographic Principle Theory of Mind MARK GERMINE Institute for Psycho science (Wikipedia).** Discordant-discrepant, dissonant-hard-line, implacable, incompatible, incongruous- inconsistent, inexorable-inflexible, inharmonious-intransigent, opposed, reluctant, unappeasable, uncompromising, unfriendly some of these expatiations and enumerations might sound, a more generalised interpretations is given by a Vedantin. Look at this arresting passage: The teacher continued to be silent. When addressed a second and third time he said: "I am teaching, but you do not follow. The Self is silence." The undetermined and unthinkable character of the Brahman is

a consequence of the absolute's eternal and immutable nature. To concede the existence of a real universe is, from the Vedantin's point of view, to posit the existence of a reality apart from the Brahman. Nor can we simply identify a real universe with the absolute unless we are prepared to compromise its unchanging, absolute status. The criterion of authenticity is immutability. Reality never changes; only that which is less than real can appear to do so. Reality is constant in the midst of change. What this means essentially is that there is change although nothing changes. This impossible situation is reflected in the ultimate impossibility of change itself. That which does not exist prior to its changing and at the end, after it has changed, must be equally non-existent between these two moments. Although the world of change appears to be real, it cannot be so. Change, according to the Vedantin, presupposes a loss of identity. Reality cannot suffer transformation; if it were to do so, it would become something else and the real would be deprived of its reality. The immortal can never become mortal, nor can the mortal become immortal. The ultimate nature of anything cannot change. Change of any sort is merely apparent (vivarta); the world of change and becoming is a false super- imposition (adhyaropa, adhyasa) on the absolute. In cosmic terms, the mistake (bhranti) consists of the supposition that the real Brahman is the unreal universe and the unreal universe is the real Brahman. In microcosmic terms, it is the mistake of falsely conceiving the body, mind or even one's personality to be the Self. In the same way as the image of a snake is falsely superimposed on a rope, similarly the universe is falsely projected onto the real substratum, the Brahman. Ignorance is not merely a personal lack of knowledge, but a cosmic principle. As such it is called "Maya," the indefinable factor (anirvacaniya) that brings this mistake in identity about. The reality status of this cosmic illusion is also indefinable: on the one hand it is not Brahman, the sole reality; on the other hand it is not absolutely non-existent like a hare's horn or the son of a barren woman. Brahman is the source of world appearances only in the sense of being their unconditioned ground or essential nature. The universe is false not because it has no nature of its own but because it does have one. Just as the illusion of a snake disappears when one sees that it is nothing but a rope, similarly cancellation (badha) of the empirically real occurs

when the absolute reality of the Brahman is realised. Thus, according to Vedanta, appearance implies the real, while the real need not imply appearance. To appear is essentially to appear in place of the real, but to be real is not necessarily to appear. All things -exist because the absolute exists. It is their Being. Thus the very existence of phenomena implies their non-existence as independent realities.

Part Each region created encompasses points with the same number of roots (counting all multiplicities)with positive real parts the region where this number is zero is the region of local stability for a particular value of $\mathcal{T}_9$ To define absolute stability it is necessary to prove that a region exists independently of $\mathcal{T}_9$

The - $D_9$ positions are defined by letting $\lambda_9 = iZ_9$ and separating (A1.2) into real and imaginary parts:

$$Z_9^2 + V_9 Z_9 \int_0^\infty \sin(Z_9 T_9) f(T_9) dT_9 - (X_9 + Y_9) = 0$$

$$Z_9 \left( U_9 - V_9 \int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9 \right) = 0$$

Noting that $X_9 + Y_9 = 0$ is we assume X+Y=0.It follows that $Z_9 = 0$ is not a solution of (A1.3) and from (A1.4) that a D- partition can only be crossed if

$$U_9 = V_9 \int_0^\infty cos\left(Z_9 T_9\right) f\left(T_9\right) dT_9$$

If $U_9 > |V_9|$then given the constraints of $f(T_9)$, $(A1.5)$*can never be satisfied Co*nsequently if a distributed delay is included in an absolutely stable.

Note that the argument is unchanged even when the distribution (including the mean) changes over time provided that the changes are random with respect to the system As long as the left-hand side of (A9.5)remains greater than the right-hand side then the system will approach its equilibrium even though the rate of approach will vary.

ABSOLUTE STABILITY WITH A DELAY DUE TO CLEDIT CONSDIDATION

It is shown elsewhere (Nunney, 1985 a) that the characteristic equation of the model (5) is given by

$$\lambda_9^2 + (U_9 - V_9 + M_9)\lambda_9 - M_9\lambda_9 e^{-T_9\lambda_9} + [(U_9 - V_9)M_9 + Y_9] - [((U_9 -$$

V 9M 9–]**e**–*T*9λ 9=0

(For definitions see Appendix9) We follow the technique outlined in Appendix 9 and let $\lambda_9 = iZ_9$ and separate (A2.9) into real and imaginary parts

$$Z_9^2 + M_9 Z_9 \sin(Z_9 T_9) - (A_9 M_9 + Y_9) + (A_9 M_9 + X_9)\cos(Z_9 T_9) = 0$$

$$(A_9 M_9) Z_9 - M_9 Z_9 \cos(Z_9 T_9) + (A_9 M_9 - X_9)\sin(Z_9 T_9) = 0$$

Quantum theory representations of real (R) and complex (C) numbers are given that is based on (e) states of single, finite strings of (e) qukits for (e) any base k > 1. Both unary representations and (e&eb) the possibility that qukits with (e&eb) k a prime number are (=) elementary and the rest composite are discussed. Cauchy sequences of (e) qukit string states are defined from (e) the arithmetic properties. The representations of (e) R and C, as equivalence classes of these sequences, differ from (e&eb) classical kit string state representations in two ways: the freedom of choice of basis states, and the fact that each quantum theory representation is (=) part of a mathematical structure that is itself based on (e) the real and complex numbers. These aspects enable (eb) the description of 3 dimensional frame fields labeled by (e) different k values, different basis or gauge choices, and different iteration stages. The reference frames in the field are based on (e) each R and C representation where (e) each frame contains (e) representations of *(e) all physical theories as mathematical structures based on (e) the R and C representation. Approaches to integrating this with physics are described. It is observed that R and C values of physical quantities, matrix elements, etc. which are viewed in (eb) a frame as elementary and featureless, are seen in a parent frame as equivalence classes of (e) Cauchy sequences of qukit string states. Subjects:Quantum Physics (quant-ph); High

Energy Physics - Theory (hep-th); General Mathematics (math.GM) Journal reference:Advances in Quantum Computation, K. Mahdavi and D. Koslover, Eds., Contemporary Mathematics Vol. 482, pp125-163, (2009) Cite as: arXiv:0709.2664 [quant-ph] (or arXiv:0709.2664v2 [quant-ph] for this version) **Reference Frame Fields based on Quantum Theory Representations of Real and Complex Numbers Paul Benioff STOP** This paper extends earlier work on quantum theory representations of natural numbers N, integers I, and rational numbers Ra to describe (eb) a space of these representations and transformations on (e&eb) the space. The space is parameterized by (e) 4-tuple points in a parameter set. Each point, (k,m,h,g), labels (e&eb) a specific representation of X = N, I, Ra as a Fock space F^{X}_{k,m,h} of (e) states of (e) finite length strings of qukits q and a string state basis B^{X}_{k,m,h,g}. The pair (m,h) locates the q string in (eb) a square integer lattice I \times I, k is the q base, and the function g fixes the gauge or basis states for (e) each q. Maps on the parameter set induce (eb) transformations on on the representation space. There are two shifts, a base change operator W_{k',k}, and (e&eb) a basis or gauge transformation function U_{k}. The invariance of (e) the axioms and theorems for N, I, and Ra under (e&eb) any transformation is discussed along with (e&eb) the dependence of the properties of W_{k',k} on (e&eb) the prime factors of k' and k. This suggests (eb) that one consider prime number q's, q_ {2}, q_{3}, q_{5}, etc. as (=) elementary and the base k q's as composites of (e) the prime number q's. Subjects: Quantum Physics (quant-ph); High Energy Physics - Theory (hep-th); Mathematical Physics (math-ph) Cite as: arXiv: 0704.3574 [quant-ph] (or arXiv: 0704.3574v1 [quant-ph] for this version) **Space of Quantum Theory Representations of Natural Numbers, Integers, and Rational Numbers Paul Benioff STOP** In this paper fields of quantum reference frames based on (e) gauge transformations of (e) rational string states are described in a way that, hopefully, makes (eb) them (eb) more understandable than their description in an earlier paper. The approach taken here is based on three main points: (1) There are (=) a large number of different quantum theory representations of (e) natural numbers, integers, and rational numbers as states of qubit strings. (2) For each representation, Cauchy sequences of rational string states give (eb) a representation of the real (and complex) numbers. A reference frame is associated to (e&eb) each representation. (3) Each frame contains (e) a representation of (e) all mathematical and physical theories that have (e) the representations of the real and complex numbers as (=) a scalar base for the theories. These points and other aspects of the resulting fields are then discussed and justified in some detail. Also two different methods of relating the frame field to (e&eb) physics are discussed. Comments: 24 pages, 5 figures; submitted to proceedings, 3rd Feynman Festival, University of Maryland, 2006 Journal reference: Journal of Physics: Conference Series 70 **Fields of quantum reference frames based on different representations of rational numbers as states of qubit strings Paul Benioff STOP System studied is as follows:** System and attendant properties thereof is studied. Gravity Kills Schrödinger's Cat , Warped Space-Time Prevents Quantum Superpositions Of Large-Scale Objects, Error Correcting Codes In Quantum Theory, Scheme For Reducing Decoherence In Quantum Computer Memory, Correcting Quantum Computer Errors, Qubits Based On Quantum Spins, Connection Between The Observers Will Typically Be Subject To Disturbances That Can Disrupt The Entanglement, Relationship Between Communication Complexity Problems Bell Nonlocal Correlations, Bell Nonlocal Correlations, Communication Complexity And The Violation Of A Bell Inequality, Exchange Of The Information After Simultaneous Teleportation To Reveal The Path Of The Teleported System In A Three-Round Protocol, Efficient Teleportation Protocol

Where $A_9 = U_9 - V_9$Using (A2.3), substitute for $\boldsymbol{Z}_9^2$ and $Z_9$in (A2.2):

$W_9^2 (X_9^2 - 3 M_9 A_9 X_9 - M_9^2 (X_9 - Y_9) + 2 M_9^3 A_9 + 2 A_9^2 M_9^2]$

$$+W\left[MA^3 - A^2 X + 2MA(2X + Y) - 2X^2\right] + A^2 (X + Y) = 0$$

*where* $W_9$=1-cos ($Z_9 T_9$).

In the absence of a delay the system is stable for $A_9 > 0$, and it can be shown (Nunney 1985)that a necessary condition for stability in the presence of an infinite delay is $A_9 > (X_9 - Y_9)/2\ M_9$ Thus a necessary condition for absolute stability is

$A_9 > [\ X_9 - Y_9/2\ M_9$

0 whichever is the larger

Using(A2.5) as a starting point we now search for constraints upon $A_9$ which exclude the possibility of (A2.4)being satisfied by a real value of $W_9$ in its permissible range of 0-2 such constraints define absolute stability

For simplicity we rewrite (A2.4) as

$$Q(W_9, A_9) = 0$$

Which is quadratic in $W$ the shape of the quadratic function in the range $0 \leq W_9 \leq 2$ *is* of interest First consider the extreme values

$Q_9(0, A_9) = A_9^2\ (X_9 - Y_9)$

$Q_9(2, A_9) = (A_9 + 2\,M_9)^2\,(2A_9 M_9 - (X_9 - Y_9))$

*which* are never negative given (A2.5)Next note that

$$\frac{\partial Q_9}{\partial A_9} = W_9^2\left(2M_9^3 + 4A_9 M_9^2 - 3M_9 X_9\right) + W_9\left(3M_9 A_9^2 - 2A_9 X_9 + 2M_9\left(2X_9 + Y_9\right)\right)$$

$$+2A_9\left(X_9+Y_9\right)$$

$$=2M_9^2W_9^2(M_9+\mathrm{A}_9)+M_9A_9^2W_9+\left(2A_9M_9-X_9+Y_9\right)(A_9W_9+M_9W_9^2$$

$$+(2-W_9)((2M_9X_9+M_9Y_9)W_9+A_9(X_9+Y_9)),$$

Which is always positive in the range of interest Consequently the limit of absolute stability occurs either when the function $Q_9$quadratic in $W_9$touches the $W$–axis only at its minimum value in which case (A2.9) shows that further increases in $A$ will life the function away from the axis and maintain absolute stability or when the function $Q_9$has no minimum in the range 0 $W_9$ 2.Defining

$Q_9\left(W_9,A_9\right)=a_9W_9^2+c_9$ *where* $a_9b_9$ *and* $c_9$ Are functions of $A_9$we can define the condition for a minimum in the range 0-2

To Heidegger writings we shall compose the model. Following is not still written model able form. For one of the models please see any one of the papers under the appellation: Ïn Excited Reveries

In his later writings Heidegger drops the term 'phenomenon' in preference for the verbal form 'phainesthai' in order to emphasize even more the actuality or presentational property of Being. Explaining this new form of the term he writes: "Being disclosed itself to the ancient Greeks as 'physis'. The etymological roots 'phy-' and 'pha-' designate the same thing: 'phyein', the rising-up or upsurge which resides within itself as 'phainesthai', lighting-up, self-showing, coming-out, appearing-forth." Heidegger contrasted his notion of phenomenon with semblance (Schein) and with appearing (Erscheinung). In the case of semblance a thing can show itself as that which it is not, as when fool's gold shows itself to be gold. The ancients always allied semblance with non-being. Heidegger points out, however, that semblances are grounded in

showings, and so does Abhinava. Both Heidegger and Abhinava consequently maintain that all semblances have a real basis and are to be treated as instances of phenomena along with the so-called real showing or manifestation of non-deceptive objects. So Heidegger states that: 'how- ever much seeming, just that much being'. Thus self-showing or appearing defines Being as phenomenon, but this definition of Being is as yet incomplete. Being is not only self-showing but 'logos' which Heidegger explains means 'discourse' (Rede) in the sense of 'apophansis': 'letting-be-seen'. Phenomenology, which according to Heidegger is the only correct study of Being, means 'letting-be-seen-that-which- shows-itself. This is true of Saiva Paramadvaita as well. The reality of the world demands recognition; we are forced to accept the direct presentation of the fact of our daily experience. As Abhinava says: "if practical life, which is useful to all persons at all times, places and conditions were not real, then there would be nothing left which could be said to be real" A thousand proofs could not make 'blue' other than the colour blue. The reality of whatever appears in consciousness cannot be denied. Objects appear; they do not cease to do so by a mere emphatic denial. The manifestation of an entity in its own specific form is a fact at one level of consciousness; it is real. The appearing of the same entity in the same form but recognised to be a direct representation of the absolute is also a fact, but at another level of consciousness. It is no more or less real than the first. 'As is the state of consciousness, so is the experience,' says Abhinava. Although the nature of the absolute is discovered at a higher level of consciousness, The Doctrine of Vibration nonetheless it presents itself to us directly in the specific form in which we perceive things; otherwise there would be no way in which we could penetrate from the level of appearing to that of its source and basis. Abhinava writes: Real is the entity iyastu) that appears in the moment of direct perception (sak$atkara), that is to say, within our experience of it. Once its own specific form has been clearly determined one should, with effort, induce it to penetrate into its pure conscious nature. All things are known to be just as they present themselves. The concrete actuality of

being known (pramiti), irrespective of content, is itself the vibrant (spanda) actuality of the absolute. Liberating knowledge is gained not by going beyond appearances but by attending closely to them. "The secret," Mahesvarananda says, "is that liberation while alive (jivanmukti) is the profound contemplation of Maya's nature." No ontological distinction can be drawn between the absolute and its manifestations because both are an appearing (dbhasa), the latter of diversity and the former of 'the true light of consciousness which is beyond Maya and is the category Siva'. Those who have attained the category of Pure Knowledge above Maya and have thus gone beyond the category of Maya, see the entire universe as the light of consciousness . . . Just as the markings [on a feather] are nothing apart from the feather, the feather [is nothing apart from] them, similarly, when the light of consciousness is manifest, the whole group of phenomena is manifest as the light of consciousness itself. Within the sphere of Maya, every entity's 'own nature' (svabhdva) corresponds to its specific manifest form. Accordingly it is defined as that which distinguishes it from all else and from which it never deviates. Above the sphere of Maya, that is, above the level of objectivity, is the domain of the subject. At this level, everything is realised to be part of the fullness of the experience r and hence no longer bound by the conditions which impinge on the object. Here the part is discovered to be the whole, that is, consciousness in toto. In this sphere beyond relative distinctions, the yogi realises that (all) the categories of existence are present in every single category. m The yogi experiences every individual particular as the sum total of everything else. He recognises that all things have one nature and that every particular is all things. This is the 'essence' (sard) or co-extensive unity (samarasya) of all things. Integral Monism of Kashmiri SaivismWe have established that reality is manifest according to how [and the degree in which] the freedom of consciousness reveals it and that [this freedom] is the womb of all forms. Just as 'sweetness' is present in its entirety in every atom of the sugarcane, so each and every atom [of the universe] bears within itself the emanation of all things. This is

the level of consciousness in which the absolute reflects on itself realising to its eternal delight and astonishment (camatkara) its own integral nature. The reality of the world of diversity is not denied, but experienced in a new mode of awareness free of time and space in the eternal omnipresence of the Here and Now. [Phenomenal forms of awareness] such as 'this [exists]', born of the colouring [imparted to the absolute] by the limitations engendered by the diversifying power of time {kalakalana) also emanate within the Supreme Principle. There [at that level], Fullness {purriata) is the one nature [of all things] and so everything is omnipresent; otherwise, associated with division (khancjlana), the Fullness [of the absolute] would not be full. The content of absolute consciousness consists of diverse appearings (abhdsa) which, because they are manifest through it in this way, do not compromise the wholeness of consciousness. Everything we perceive is a momentary collocation of a number of such manifestations which combine together like 4 a row of altar lamps' (dipavalT) to form the single radiant picture of the universe. The individual objects which constitute the universe are specific collocations of such 'atomic' appearings. Together they form a single unified particular which appears according to its own defining features (svalaksand). A jar, for example, consists of a number of appearances such as 'round', 4 fat\ 'earthen', 'red', etc., which together discharge a single function (arthakriya), in this case, that of carrying the appearance 'water'. They unite with each other much as the scattered rays of a lamp come together when focused, or as the various currents of the sea together give rise to waves. Atomic appearings can combine in any number of ways, provided that they are not contrary to one another as established by the dictates of natural law (niyati). An appearance of 'form', for example, cannot combine with that of 'air'. Insofar as they share a common basis (samanyadhikaranya), a given cluster of appearances appears as a single whole. This common basis is the most prominent member of the group; the appearance 'jar' is such in the example quoted above. Any one appearance in a cluster may assume a more important or subordinate role.

And the condition for a real solution to (A2.6) in that range if both condition are satisfied then system cannot be absolutely stable if follows that given (A2.5) a system will not be absolutely stable if

$$4a_9 - b_9 < 0$$

And

$$b_9^2 > 4a_9c_9$$

$$Case(i): Y_9 = 0$$

$Let\ Y_9 = 0$ And try the substitution A= $X_9 = /2M_9\ in$ (A2.4) this leads directly to a linear equation with the solution

$W_9 = 1$The quadratic conditions (A2.10) and (A2.10) and (A2.11) do not apply: however note that $\boldsymbol{Q}$ ($W_9 X_9 /$ $2M_9$) is positive for *W<2. Since* $Q_9$ increases with $A_9$(see A2.9), increasing $A_9$will generate no further solutions of $Q_9$=0. Hence the condition A>X/2M defines absolute stability when $Y_9 = 0$.

$Case\ (ii): Y_9 = y_9\ X_9$ ,

$Let\ Y_9 = y_9\ X_9$And try the substitution $A_9 = X_9/(2 + A_9\ y_9)M$ in (A2.4), where n is any real number defining A>0 after rearrangement and dividing through by $X_9/(2 + n_9 y_9)$ we have

$$W_9^2\ y_9\left[X_9 n_9 \frac{(1 + n_9 y_9)}{(2 + n_9 y_9)} + M_9^2\left(2 - n_9 + n_9 y_9\right)\right]$$

$$-W_9\,X_9\left[\frac{X_9\left(1+n_9y_9\right)}{M_9^2\left(2+n_9y_9\right)^2}+2\,y_9\left(n_9-1\right)+\frac{X_9^2\left(9+y_9\right)}{M_9^2\left(2+n_9y_9\right)}\right]=0.\ (A2.12)$$

Since $n_9$ and $A_9$ are interchangeable we need only to determine the conditions for absolute stability in terms of $n_9$

Defining $E_9=\left(2+n_9y_9\right),$ *the condition*$(A2.99?)$*can be written as*

$$\left[\frac{X_9\left(E_9-9\right)}{E_9^2}\right]^2\left[\frac{X_9^2}{M_9^2}-4\,y_9\,E_9^2\frac{X_9}{M_9^2}+4y_9\left(n_9-2\right)E_9^3\right]\ 0.$$

Reversal of this inequality gives a sufficient condition for absolute stability [no real solution to (A2.6)] which upon solving the quadratic in $(X_9/W_9^2)$ can be expressed as

$$2\,E_9\left[y_9E_9+\left(y_9\,E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}\right]>\frac{X_9}{M_9^2}>2E_9\left[y_9E_9-y_9E_9\left(9+y_9\right)\left(E_9-n_9\right)\right)^{9/2}]$$

dynamic service migration problem in mobile edge-clouds that host cloud-based services at the network edge. This offers the benefits of reduction in network overhead and latency but requires service migrations as user locations change over time. It is challenging to make these decisions in an optimal manner because of the uncertainty in node mobility as well as possible non-linearity of the migration and transmission costs. In this paper, we formulate a sequential decision making problem for service migration using the framework of Markov Decision Process (MDP). Our formulation captures general cost models and provides a mathematical framework to design optimal service migration policies. In order to overcome (e) the complexity associated with (e&eb) computing the optimal policy, we approximate the underlying state space by the distance between the user and service locations. We show that the resulting MDP is exact for uniform one-dimensional mobility while it provides a close approximation for uniform two-dimensional mobility with a constant additive error term. We also propose a new algorithm and a numerical technique

for computing the optimal solution which is significantly faster in computation than traditional methods based on value or policy iteration. We illustrate the effectiveness of our approach by simulation using real-world mobility traces of taxis in San Francisco. Comments: in Proc. of IFIP Networking 2015Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1506.05261 [cs.DC] (or arXiv: 1506.05261v1 [cs.DC] for this version) **Dynamic Service Migration in Mobile Edge-Clouds Shiqiang Wang, Rahul Urgaonkar, Murtaza Zafer, Ting He, Kevin Chan, Kin K. Leung STOP** Let kE denote the group algebra of an elementary abelian p-group of rank r over an algebraically closed field of characteristic p. We investigate the functors Fi from kE-modules of constant Jordan type to vector bundles on Pr−1(k), constructed by Benson and Pevtsova. For a kE-module M of constant Jordan type, we show that restricting the sheaf Fi(M) to a dimension s−1 linear subvariety of Pr−1(k) is equivalent to restricting M along a corresponding rank s shifted subgroup of kE and then applying Fi. In the case r=2, we examine the generic kernel filtration of M in order to show that Fi(M) may be computed on certain sub quotients of M whose Loewy lengths are bounded in terms of i. More precise information is obtained by applying similar techniques to the nth power generic kernel filtration of M. The latter approach also allows us to generalise our results to higher ranks r. Subjects: Representation Theory (math.RT) MSC classes: 16G10, 20C20 Cite as: arXiv: 1504.01994 [math.RT] (or arXiv: 1504.01994v1 [math.RT] for this version) **Modules of constant Jordan type, pullbacks of bundles and generic kernel filtrations Shawn Baland, Kenneth Chan STOP** Mobile micro-cloud is an emerging technology in distributed computing, which is aimed at providing seamless computing/data access to the edge of the network when a centralized service may suffer from poor connectivity and long latency. Different from the traditional cloud, a mobile micro-cloud is smaller and deployed closer to users, typically attached to a

cellular basestation or wireless network access point. Due to (e) the relatively small coverage area of each basestation or access point, when a user moves across areas covered by different base stations or access points which are attached to (e&eb) different micro-clouds, issues of service performance and service migration become (=) important. In this paper, authors consider such migration issues. They model the general problem as (=) a Markov decision process (MDP), and show that, in the special case where (e) the mobile user follows(eb) a one-dimensional asymmetric random walk mobility model, the optimal policy for (e) service migration is (=) a threshold policy. Authors obtain the analytical solution for (e) the cost resulting from arbitrary thresholds, and then propose an algorithm for finding the optimal thresholds. The proposed algorithm is more efficient than standard mechanisms for solving MDPs. Comments: in Proc. of IEEE MILCOM 2014, Oct. 2014 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) DOI: 10.1109/MILCOM.2014.145 Cite as: arXiv: 1503.05141 [cs.DC] (or arXiv: 1503.05141v1 [cs.DC] for this version) **Mobility-Induced Service Migration in Mobile Micro-Clouds Shiqiang Wang, Rahul Urgaonkar, Ting He, Murtaza Zafer, Kevin Chan, Kin K. Leung STOP** Seamless computing and data access is enabled by the emerging technology of mobile micro-clouds (MMCs). Different from traditional centralized clouds, an MMC is typically connected directly to a wireless base-station and provides services to a small group of users, which allows users to have instantaneous access to cloud services. Due to the limited coverage area of base-stations and the dynamic nature of mobile users, network background traffic, etc., the question of where to place the services to cope with these dynamics arises. In this paper, we focus on dynamic service placement for MMCs. We consider the case where there is an underlying mechanism to predict the future costs of service hosting and migration, and the prediction error is assumed to be bounded. Our goal is to find the optimal service

placement sequence which minimizes the average cost over a given time. To solve this problem, we first propose a method which solves for the optimal placement sequence for a specific look-ahead time-window, based on the predicted costs in this time-window. We show that this problem is equivalent to a shortest-path problem and propose an algorithm with polynomial time-complexity to find its solution. Then, we propose a method to find the optimal look-ahead window size, which minimizes an upper bound of the average cost. Finally, we evaluate the effectiveness of the proposed approach by simulations with real-world user-mobility traces. Comments: in Proc. of IEEE ICC 2015 Subjects: Distributed, Parallel, and Cluster Computing (cs.DC); Networking and Internet Architecture (cs.NI); Optimization and Control (math.OC) Cite as: arXiv: 1503.02735 [cs.DC] (or arXiv: 1503.02735v1 [cs.DC] for this version) **Dynamic Service Placement for Mobile Micro-Clouds with Predicted Future Costs Shiqiang Wang, Rahul Urgaonkar, Kevin Chan, Ting He, Murtaza Zafer, Kin K. Leung STOP** Authors give a mathematically rigorous construction of Rozansky-Witten's 3-dimensional σ-model as a perturbative quantum field theory (QFT) by applying Costello's approach using the Batalin-Vilkovisky (BV) formalism. The quantization of our model is obtained via the technique of configuration spaces. We also investigate the observable theory following the work of Costello-Gwilliam. In particular, we show that the cohomology of local quantum observables on a genus g handle body is given by $H*(X,(\wedge*TX)\otimes g)$, where X is the target hyperk\"ahler manifold. We further give a mathematical definition of the partition function and prove that it coincides with the Rozansky-Witten invariants. Subjects: Quantum Algebra (math.QA); High Energy Physics - Theory (hep-th); Algebraic Geometry (math.AG); Differential Geometry (math.DG); Geometric Topology (math.GT) Cite as: arXiv:1502.03510 [math.QA] (or arXiv:1502.03510v1 [math.QA] for this version) **A mathematical foundation of Rozansky-Witten theory Kwokwai Chan, Naichung Conan Leung, Qin Li STOP Shiraz Minwalla1, Mark Van**

**Raamsdonk1 and Nathan Seiberg2** study the perturbative dynamics of (e) noncommutative field theories on (eb) Script Rd, and find (eb) an intriguing mixing of the UV and the IR. High energies of virtual particles in loops produce (eb) non-analyticity at low momentum. Consequently, the low energy effective action is (=) singular at zero momentum even when (e) the original noncommutative field theory is (=) massive. Some of the nonplanar diagrams of (e) these theories are (=) divergent, but they interpret these divergences as (=) IR divergences and deal with (e&eb) them accordingly. Authors explain how this UV/IR mixing arises from (e) the underlying noncommutativity. This phenomenon is reminiscent of (e) the channel duality of (e) the double twist diagram in (eb) open string theory. **Noncommutative perturbative dynamics Shiraz Minwalla1, Mark Van Raamsdonk1 and Nathan Seiberg2 Published 15 March 2000 • Journal of High Energy Physics, Volume 2000, JHEP02 (2000 STOP** Authors give a derivation of the gravitational Hamiltonian starting from (e) the Einstein - Hilbert action, keeping track of (e) all surface terms. This derivation can be applied to (e&eb) any spacetime that asymptotically approaches (eb) a static background solution. The surface term that arises in (eb) the Hamiltonian can be taken as (=) the definition of the `total energy', even for (e) spacetimes that are not (e) asymptotically flat. (In the asymptotically flat case, it agrees with (eb) the usual ADM energy.) They also discuss the relation between the Euclidean action and (e&eb) the Hamiltonian when (e) there are horizons of infinite area (e.g. acceleration horizons) as well as (e&eb) the usual finite area black hole horizons. Acceleration horizons seem to be (=) more analogous to (e) extreme than nonextreme black holes, since they find evidence that their horizon area is not (e) related to the total entropy **The gravitational Hamiltonian, action, entropy and surface terms S W Hawking and Gary T Horowitz Classical and Quantum Gravity, Volume 13, Number 6 http://dx.doi.org/10.1088/0264-9381/13/6/017 STOP** Recent observations of Type 1a supernovae indicating (eb) an accelerating universe have (e) once more drawn attention to the possible existence, at the

present epoch, of (e) a small positive Λ-term (cosmological constant). In this paper authors review both observational and theoretical aspects of (e) a small cosmological Λ-term. Authors discuss the current observational situation focusing on cosmological tests of Λ including (e) the age of the universe, high redshift supernovae (e&eb) gravitational lensing (e&eb) galaxy clustering and (e&eb) the cosmic microwave background. They also review the theoretical debate surrounding Λ: the generation of Λ in models with (e&eb) spontaneous symmetry breaking and through(e&eb) quantum vacuum polarization effects — mechanisms which are known to give rise to (eb) a large value of Λ hence leading to (eb) the "cosmological constant problem." More recent attempts to generate (eb) a small cosmological constant at the present epoch using (e) either field theoretic techniques or by (e) modelling a dynamical Λ-term by scalar fields are also extensively discussed. Anthropic arguments favouring a small Λ-term are briefly reviewed. A comprehensive bibliography of recent work on Λ is provided. **ARUN SAHNI and ALEXEI STAROBINSKY, Int. J. Mod Phys. D 09, 373 (2000) DOI: 10.1142/S0218271800000542 THE CASE FOR A POSITIVE COSMOLOGICAL Λ-TERM STOP** Authors explore the observational adequacy of (e) a class of unified dark energy and dark matter (UDE/M) models with (e&eb) a fast transition. Constraints are set using (e) a combination of geometric probes—some low redshift ones and (e&eb) some high redshift ones (CMB-related included). The transition is phenomenologically modeled by (e) two different transition functions corresponding to (e&eb) a fast and to an ultrafast transition, respectively. They find that the key parameters governing (e&eb) the transition can be (=) well constrained, and from (e) the statistical point of view it follows (eb) that the models cannot be discarded when compared to (e&eb) LCDM. Authors find the intriguing result that standard or input parameters such as Ωm and Ωb are (=) far better constrained than in (eb) LCDM, and this is the case for (e) the derived or output parameter measuring (eb) the deceleration value at present, q0. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043537

© 2016 American Physical Society **Cosmological constraints on fast transition unified dark energy and dark matter models Ruth Lazkoz, Iker Leanizbarrutia, and Vincenzo Salzano Phys. Rev. D 93, 043537 – Published 22 February 2016 STOP** Authors report on constraints to (e) the cosmological jerk parameter (j) and to possible variability of (e) the fine-structure constant ($\Delta\alpha/\alpha$) based on (e) stochastic quintessence models of (e) dark energy, discussed by Chongchitnan and Efstathiou [Phys. Rev. D 76, 043508 (2007)]. Authors confirm the results by (e) these authors in the sense that many viable solutions can be obtained, obeying (e&eb) current observational constraints in low redshifts. They add the observables j and $\Delta\alpha/\alpha$ to (e&eb) this conclusion. However, they find peculiarities that may produce (eb) in the nearby universe, potential observational imprints in (eb) future cosmological data. Authors conclude, for redshifts $z\lesssim 3$, that (i) j(z) fluctuates due to (e&eb) the stochasticity of the models, reaching (eb) an amplitude of up to 5% relatively to (e&eb) the $\Lambda$ cold dark matter model value ($j\Lambda CDM=1$); and (ii) by contrasting two distinct ("extreme") types of solutions, variabilities in $\alpha(z)$, linked to (e&eb) a linear coupling ($\zeta$) between the dark energy and (e&eb) electromagnetic sectors, are (eb) weakly dependent on redshift, for (e) couplings of the order $|\zeta|\sim 10-4$, even for (e) large variations in the equation of state parameter at (eb) relatively low redshifts. Nonlinear couplings produce (eb) an earlier and steeper onset of (e) the evolution in $\Delta\alpha/\alpha$ (z), but can still accommodate (e) the data for weak enough couplings. DOI: http://dx.doi.org/10.1103/PhysRevD.93.043509 © 2016 American Physical Society **The Void Galaxy Survey: Star Formation Properties B. Beygu1,6,★, K. Kreckel2, J. M. van der Hulst1, T. H Jarrett3, R. Peletier1, R. van de Weygaert1, J. H. van Gorkom4 and M. A. Aragon-Calvo5 STOP Stochastic quintessence models: Jerk and fine-structure variability constraints Christine C. Dantas and André L. B. Ribeiro Phys. Rev. D 93, 043509 – Published 8 February 2016 STOP** Authors study the star formation properties of 59 void galaxies as part of the Void Galaxy Survey (VGS). Current star formation

rates are derived from (e) $\rm {H\alpha} $ and recent star formation rates from (e) near-UV imaging. In addition, infrared 3.4 $\rm {\mu m}$, 4.6 $\rm {\mu m}$, 12 $\rm {\mu m}$ and 22 $\rm {\mu m}$ WISE emission is used as (=) star formation and mass indicator. Infrared and optical colours show that (eb) the VGS sample displays (eb) a wide range of dust and metallicity properties. They combine these measurements with (e&eb) stellar and $\rm {H\small {I}} $ masses to measure (eb) the specific SFRs ($\rm {SFR/M_ {*}} $) and (e&eb) star formation efficiencies ($\rm {SFR/M_ {H\small {I}}}$). Authors compare the star formation properties of their sample with (e&eb) galaxies in (eb) the more moderate density regions of (e) the cosmic web, 'the field'. They find that specific SFRs of (e) the VGS galaxies as (=) a function of stellar and $\rm {H\small {I}}$ mass are similar to (e&eb, =) those of the galaxies in these field regions. Their $\rm {SFR\alpha} $ is slightly elevated than (e) the galaxies in the field for (e) a given total $\rm {H\small {I}} $ mass. In the global star formation picture presented by (e) Kennicutt-Schmidt, VGS galaxies fall into (eb)the regime of low average star formation and correspondingly (e&eb) low $\rm {H\small {I}} $ surface density. Their mean $\rm {SFR\alpha /M_{H\small {I}}}$ and (e&eb) $\rm {SFR\alpha /M_{*}}$ are of (e) the order of $\rm {10^{-9.9}}$ $\rm {yr^{-1}}$. Authors conclude that while the large scale underdense environment must play (e&eb) some role in galaxy formation and growth through accretion, they find that even with respect to other galaxies in the more mildly underdense regions, the increase in star formation rate is (=) only marginal. Key words galaxies: star formation galaxies: formation —galaxies: structure large-scale structure of universe© 2016 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society Authors study the formation and evolution of (e) filamentary configurations of dark matter haloes in (eb) voids. Extant investigation uses (e) the high-resolution Λ cold dark matter simulation Cosmo Grid to (e) look for (e) void systems resembling the VGS_31 elongated system of (e) three interacting galaxies that was recently discovered by (e) the Void

Galaxy Survey inside (eb) a large void in the Sloan Digital Sky Survey galaxy redshift survey. H I data revealed these galaxies to be embedded in (eb) a common elongated envelope, possibly embedded in (eb) intravoid filament. In the Cosmo Grid simulation authors look for systems similar to (e) VGS_31 in mass, size and environment. They find (eb) a total of eight such systems. For these systems, authors study the distribution of (e) neighbour haloes, the assembly and evolution of (e) the main haloes and (e&eb) the dynamical evolution of the haloes, as well as the evolution of (e) the large-scale structure in which the systems are (=) embedded. The spatial distribution of the haloes follows that of (e&eb) the dark matter environment. Authors find that VGS_31-like systems have (e) a large variation in formation time, having (e) formed between (e&eb) 10 Gyr ago and the present epoch. However, the environments in which the systems are embedded evolved to (e) resemble each other substantially. Each of the VGS_31-like systems is embedded in (eb) an intravoid wall, that no later than z = 0.5 became the only prominent feature in its environment. While parts of the void walls retain (eb) a rather featureless character, authors find that around half of them are marked by (e) a pronounced and rapidly evolving substructure. Five haloes find themselves in a tenuous filament of a few h−1 Mpc long inside (eb) the intravoid wall. Finally, authors compare the results to (e&eb) observed data from VGS_31. Study implies that (eb) the VGS_31 galaxies formed in (eb) the same (proto) filament, and did not (e) meet just recently. The diversity amongst the simulated halo systems indicates (eb) that VGS_31 may not (e) be typical for (e) groups of galaxies in voids. Key words galaxies: formation galaxies: interactions cosmology (e&eb) theory dark matter large-scale structure of Universe © 2013 The Authors Published by Oxford University Press on behalf of the Royal Astronomical Society **Assembly of filamentary void galaxy configurations Steven Rieder1, 2, Rien van de Weygaert3, Marius Cautun3, Burcu Beygu3 and Simon Portegies Zwart1First published online August 20, 2013 STOP** Authors measure the star formation properties of (e) two large samples of

galaxies from (e) the SDSS in large-scale cosmic voids on (e&eb) time scales of 10 Myr and 100 Myr, using (e) Hα emission line strengths and GALEX FUV fluxes, respectively. The first sample consists of 109,818 optically selected galaxies. They find that void galaxies in this sample have (e) higher specific star formation rates (SSFRs; star formation rates per unit stellar mass) than (e&eb) similar stellar mass galaxies in (eb) denser regions. The second sample is (=) a subset of (e) the optically selected sample containing (e) 8070 galaxies with (e&eb) reliable HI detections from ALFALFA. For the full HI detected sample, SSFRs do not (e) vary systematically with (e&eb) large-scale environment. However, investigating only the HI detected dwarf galaxies reveals (eb) a trend towards higher SSFRs in (eb) voids. Furthermore, authors estimate the **star formation** rate per unit HI mass (known as the star formation efficiency; SFE) of a galaxy, as (=) a **function of environment**. For the overall HI detected population, authors notice no (e) environmental dependence. (AC=DC=0)Limiting the sample to dwarf galaxies again reveals (eb) a trend towards higher SFEs in (eb) voids. These results suggest (eb) that void environments provide (eb) a nurturing environment for (e) dwarf galaxy evolution allowing (eb) for higher specific star formation rates and efficiencies. Comments: submitted to ApJ 12-2015 Subjects: Astrophysics of Galaxies (astro-ph.GA) Cite as: arXiv: 1601.04092 [astro-ph.GA]**.** Can Grover's algorithm speed up (eb+) search of a physical region—for example a 2-D grid of size $\sqrt{n} \times \sqrt{n}$? The problem is that $\sqrt{n}$ time seems to be needed for (e) each query, just to move amplitude across the grid. Here authors show that this problem can be surmounted (e) refuting a claim to (e) the contrary by Benioff. In particular, authors show how to search a d-dimensional hypercube in (eb) time $O(\sqrt{n})$ for $d \geq 3$, or $O(\sqrt{n}\log 5/2n)$ for $d = 2$. More generally, authors introduce a model of (e) quantum query complexity on graphs, motivated by (e) fundamental physical limits on (e&eb) information storage, particularly the holographic principle from (e) black hole thermodynamics. Results in this model include (e) almost-tight upper and lower bounds for (e) many search

tasks; a generalized algorithm that works for (e) any graph with good expansion properties, not just hypercubes; and relationships among several notions of 'locality' for (e) unitary matrices acting on graphs. As an application of results, authors give an O ( √ n)-qubit communication protocol for (e) the disjointness problem, which improves (eb+) an upper bound of Høyer and de Wolf and matches (e&eb) a lower bound of Razborov. ACM Classification: F.1.2, F.1.3 AMS Classification: 81P68, 68Q10 Key words and phrases: Quantum computing, (e&eb) Grover search, amplitude amplification, (e&eb) quantum communication complexity, disjointness, lower bounds ∗This work was mostly done while the author was a **PhD student at UC Berkeley, supported by an NSF Graduate Fellowship and by ARO grant DAAD19-03-1-0082.†Supported by an IQC University Professorship and by CIAR**. STOP Grover's quantum search algorithm provides (eb) a way to speed up (eb) combinatorial search, but is not (e) directly applicable to searching (e&eb) a physical database. Nevertheless, Aaronson and Ambainis showed (eb) that a database of N items laid out in (eb) d spatial dimensions can be searched in time of (e) order N−−√ for d>2, and in time of order N−−√poly (logN) for d=2. Authors consider an alternative search algorithm based on (e) a continuous-time quantum walk on (e&eb) a graph. The case of the complete graph gives (eb) the continuous-time search algorithm of (e) Farhi and Gutmann, and other previously known results can be used to show that (eb) N−−√ speedup can also be achieved on (e&eb) the hypercube. Authors show that full N−−√ speedup can be achieved on (eb) a d-dimensional periodic lattice for d>4. In d=4, the quantum walk search algorithm takes time of (e) order N−−√poly (logN), and (e&eb) in d<4, the algorithm does not (e) provide substantial speedup. Received 16 June 2003 DOI:http://dx.doi.org/10.1103/PhysRevA.70.022314 **Spatial search by quantum walk Andrew M. Childs and Jeffrey Goldstone Phys. Rev. A 70, 022314 – Published 23 August 2004 STOP** Quantum walks are considered in (eb) a one-dimensional random medium characterized by (e) static or dynamic disorder. Quantum interference for (e) static disorder can

lead to (eb) Anderson localization which completely hinders (e) the quantum walk and it is contrasted with (e&eb) the decoherence effect of (e) dynamic disorder having (e) strength W, where (e) a quantum to (e&eb) classical crossover at time tc∝W−2 transforms (e&eb) the quantum walk into (e&eb) an ordinary random walk with (e&eb) diffusive spreading. Authors demonstrate this localization and decoherence phenomena in (eb) quantum carpets of the observed time evolution, authors relate results to (e&eb) previously studied models of (e) decoherence for quantum walks, and examine (e&eb) in detail a dimer lattice which corresponds to(e&eb) a single qubit subject to randomness. Received 11 August 2007 DOI:http://dx.doi.org/10.1103/PhysRevA.77.022302 **Quantum walks on a random environment Yue Yin, D. E. Katsanos, and S. N. Evangelou Phys. Rev. A 77, 022302 – Published 5 February 2008 STOP** Feedback quantum control, where (e) the sample `teaches' a computer-controlled arbitrary light form generator to find (e&eb) the optimal light field, is experimentally demonstrated for (e) a molecular system. Femtosecond pulses tailored by a computer-controlled acousto-optic pulse shaper excite (eb) fluorescence from laser dye molecules in (eb) solution. Fluorescence and laser power are (=) monitored, and the computer uses (e) the experimental data and a genetic algorithm to (e) optimize population transfer from ground to (e&eb) first excited state. Both efficiency (the ratio of excited state population to (e&eb) laser energy) and (e&eb) effectiveness (total excited state population) are optimized. Potential use as an `automated theory tester' is discussed. Copyright © 1997 Elsevier Science B.V. All rights reserved. **Chemical Physics Letters Volume 280, Issues 1–2, 28 November 1997, Pages 151–158 Feedback quantum control of molecular electronic population transfer Christopher J Bardeena, Vladislav V Yakovleva, Kent R Wilsona, Scott D Carpenterb, Peter M Weberc, Warren S Warrend STOP** Quantum robots are described as (=) mobile quantum computers and ancillary systems that move in and interact with (e&eb) arbitrary environments. Their dynamics is given as (=) tasks which consist of (e)

sequences of (e) alternating computation and (e&eb) action phases. A task example is considered in which a quantum robot searches (e&eb) a space region to find the location of a system. The possibility that the search can be (=) more efficient than a classical search is examined by (e) considering use of (e) Grover's Algorithm to process the search results. For reversible searches this is (=) problematic for two reasons. One is (=) the removal of entanglements generated by (e) the search process. The other is that even if the entanglement can be avoided, (eb) the search process in 2-D space regions is (=) no more efficient than a classical search. However, quantum searches of space regions with (e&eb) 3 or more dimensions are (=) more efficient than classical searches. Reasons why quantum robots are interesting independent of these results are briefly summarized. Subjects: Quantum Physics (quant-ph) Journal reference: AMS Contemporary Math Series, Vol 305, (2002) Cite as: arXiv: quant-ph/0003006 (or arXiv: quant-ph/0003006v2 for this version) **Space Searches with a Quantum Robot Paul Benioff (Physics Division, Argonne National Lab, Argonne, IL) (Submitted on 1 Mar 2000 (v1), last revised 26 Jun 2001 (this version, v2))STOP** The relationship between the foundations of mathematics and (e&eb) physics is a topic of of much interest. This paper continues this exploration by (e) examination of the effect of space and time dependent number scaling on (e&eb) theoretical descriptions of (e) some physical and geometric quantities. Fiber bundles provide (eb) a good framework to introduce (eb) a space and time or space time dependent number scaling field. The effect of the scaling field on (e&eb) a few nonlocal physical and geometric quantities are described. The effect on gauge theories is to introduce (eb) a new complex scalar field into (e&eb) the derivatives appearing in Lagrangians. U (1) invariance of Lagrangian terms does not (e) affect the real part of the scaling field. For this field, any mass is possible. The scaling field is also shown to affect (e&eb) quantum wave packets and (e&eb) path lengths, and (e&eb) geodesic equations even on (e&eb) flat space. Scalar fields described so far in physics, are (=) possible candidates for (e) the scaling

field. The lack of direct evidence for (e) the field in physics restricts the scaling field in that the gradient of (e) the field must be close to (e) zero in (eb) a local region of (e) cosmological space and time. There are no restrictions outside (e) the region. It is also seen that the scaling field does not (e) affect comparisons of computation or measurements outputs with (e&eb) one another. However it does affect (e&eb) the assignment of numerical values to the outputs of computations or measurements. These are needed because (e) theory predictions are in terms of (e&eb) numerical values. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum Information Processing, 15(3), 1081-1102, (2016) DOI: 10.1007/s11128-015-1084-z Cite as: arXiv: 1508.01732 [math-ph] (or arXiv: 1508.01732v1 [math-ph] for this version) **Space and time dependent scaling of numbers in mathematical structures: Effects on physical and geometric quantities Paul Benioff STOP** The purpose of this paper is to put the description of number scaling and its effects on (e&eb) physics and geometry on (e&eb) a firmer foundation, and to make it more understandable. A main point is that two different concepts, number and (e&eb) number value are combined in the usual representations of (e) number structures. This is valid as (e) long as just one structure of each number type is being considered. It is not valid when (e) different structures of each number type are (=) being considered. Elements of base sets of (e) number structures, considered by (e) themselves, have no meaning. They acquire meaning or value as (e) elements of a number structure. Fiber bundles over (e&eb) a space or space time manifold, M, are described. The fiber consists of (e) a collection of many real or complex number structures and (e&eb) vector space structures. The structures are parameterized by (e) a real or complex scaling factor, s. A vector space at a fiber level, s has (e) as scalars, real or complex number structures at (eb) the same level. Connections are described that relate scalar and (e&eb) vector space structures at both neighbor M locations and (e&eb) at neighbor scaling levels. Scalar

and vector structure valued fields are (=) described and covariant derivatives of these fields are (=) obtained. Two complex vector fields, each with (e&eb) one real and one imaginary field, appear, with one complex field associated with (e&eb) positions in M and the other with (e&eb) position dependent scaling factors. A derivation of (e) the covariant derivative for (e) scalar and vector valued fields gives (eb) the same vector fields. The derivation shows (eb) that the complex vector field associated with (e&eb) scaling fiber levels is (=) the gradient of a complex scalar field. Use of these results in gauge theory shows (eb) that the imaginary part of the vector field associated with (e&eb) M positions acts (e&eb) like the electromagnetic field. The physical relevance of the other three fields, if any, is (=) not known. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Proceedings SPIE STA conference, Vol. 9500, 9500 0X, 2015 DOI: 10.1117/12.2176080 Cite as: arXiv: 1503.05600 [math-ph] (or arXiv: 1503.05600v1 [math-ph] for this version) **Principal fiber bundle description of number scaling for scalars and vectors: Application to gauge theory Paul Benioff STOP** This work uses (e) fiber bundles as a framework to describe some effects of number scaling on (e&eb) gauge theory and some geometric quantities. A description of number scaling and fiber bundles over (e&eb) a flat space time manifold, M, is followed by (e&eb) a description of gauge theory. A fiber at point x of M contains (e) a pair of scaled complex number and (e&eb) vector space structures, Ccx×Vcx for each c in GL(1,C). A space time dependent scalar field, g, determines, (eb) for each x, the scaling value of (e) the vector space structures that contain (e) the values of a vector valued matter field at (eb) x. The vertical components of (e) connections between neighboring fibers are taken to be (=) the gradient field A(x)+iB(x), of g. Abelian gauge theory for (e) these fields gives (eb) the result that B is (=) massless and no mass restrictions for (e) A. Addition of an electromagnetic field dies not (e) change these results. In the Mexican hat Higgs

mechanism B combines with (e&eb) a Goldstone boson to create massive vector bosons, the photon field, and the Higgs field. For geometric quantities the fiber bundle is (=) a tangent bundle with (e&eb) a fiber at point x containing (e) scaled pairs, Rrx×Trx of real number and (e&eb) tangent space structures for (e) each x and and nonnegative real r. B is zero everywhere. The A field affects (e&eb) path lengths and the proper times of (e) clocks along paths. It also appears in (eb) the geodesic equation. The lack of physical evidence for (e) the gradient field, A(x)+iB(x) means (eb) that it either couples very weakly to (e&eb) matter fields or that it is close to zero for all x in (eb) a local region of cosmological space and (e&eb) time. It says nothing about (e&eb) the values outside the local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th); Quantum Physics (quant-ph) Journal reference: Quantum Studies: Mathematics and Foundations, 2, 289-313, 2015 DOI: 10.1007/s40509-015-0049-3 Cite as: arXiv: 1412.1493 [math-ph] (or arXiv: 1412.1493v2 [math-ph] for this version) **Fiber bundle description of number scaling in gauge theory and geometry Paul Benioff STOP** Stronger foundations for (e) earlier work on the effects of (e) number scaling, and local mathematics is described. Emphasis is placed on the effects of scaling on (e&eb) coordinate systems. Effects of scaling are represented by (e) a scalar field, θ, which appears (eb) in gauge theories as (=) a spin zero boson. Gauge theory considerations led to (eb) the concept of local mathematics, as expressed through (e&eb) the use of universes, as collections of (e) local mathematical systems at (eb) each point, x, of a space time manifold, M. Both local and global coordinate charts are (=) described. These map M into (e&eb) either local or global coordinate systems within (eb) a universe or between (e&eb) universes, respectively. The lifting of global expressions of (e) nonlocal physical quantities, expressed by (e) space and or time integrals or derivatives on M, to (e&eb) integrals or derivatives on coordinate systems, is (=) described. The assumption of local mathematics and (e&eb) universes makes integrals and

derivatives, on (e&eb) M or on global charts, meaningless. They acquire (e) meaning only when mapped into (e&eb) a local universe. The effect of scaling, by including (e) the effect of θ into (e&eb) the local maps, is described. The lack of experimental evidence for (e) θ so far shows that (eb) the coupling constant of θ to (e&eb) matter fields must be very small compared to **the fine structure constant**. Also the gradient of θ must be (=) very small in the local region of (e) cosmological space and time occupied by (e) us as observers. So far, there are no known restrictions on (e&eb) θ or its gradient in regions of space and/or time that are far away from (e) our local region. Subjects: Mathematical Physics (math-ph); General Relativity and Quantum Cosmology (gr-qc); Quantum Physics (quant-ph) Journal reference: Quantum information and computation XII, proceedings of SPIE conference, 8-9 May, 2014, Vol. 9123, 91230Q-1 DOI: 10.1117/12.2049705 Cite as:arXiv: 1405.3217 [math-ph] (or arXiv: 1405.3217v1 [math-ph] for this version) **Effects of mathematical locality and number scaling on coordinate chart use Paul Benioff STOP** Effects of local availability of (e) mathematics (LAM) and (e&eb) space time dependent number scaling on (e&eb) physics and, especially, geometry are described. LAM assumes (eb) separate mathematical systems as (=) structures at each space time point. Extension of gauge theories to include freedom of (e) choice of scaling for number structures, and other structures based on (e) numbers, results (eb) in a space time dependent scaling factor based on (e) a scalar boson field. Scaling has no effect on (e&eb) comparison of experimental results with (e&eb) one another or with theory computations. With LAM all theory expressions are elements of (e) mathematics at some reference point. Changing the reference point introduces (eb) (external) scaling. Theory expressions with integrals or derivatives over (e&eb) space or time include scaling factors (internal scaling) that cannot be removed by (e) reference point change. Line elements and path lengths, as integrals over (e&eb) space and/or time show (eb) the effect of scaling on geometry. In one example, the scaling factor goes to 0 as the (e&eb) time goes to 0, the

big bang time. All path lengths, and values of physical quantities, are crushed to (eb) 0 as t goes to 0. Other examples have (e) spherically symmetric scaling factors about some point, x. In one type, a black scaling hole, the scaling factor goes to (e&eb) infinity as the distance, d, between any point y and x goes to 0. For scaling white holes, the scaling factor goes to (e&eb) 0 as d (e&eb) goes to 0. For black scaling holes, path lengths from (e) a reference point, z, to y become (=) infinite as (e) y approaches x. For white holes, path lengths approach (e&eb) a value much less than the unscaled distance from z to (e) x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.2014360 Cite as: arXiv: 1306.4613 [quant-ph] (or arXiv: 1306.4613v1 [quant-ph] for this version) **Effects of gauge theory based number scaling on geometry Paul Benioff STOP** In gauge theories, separate vector spaces, Vx, are assigned to (e) each space time point x with (e&eb) unitary operators as maps between basis vectors in neighboring Vx. Here gauge theories are extended by replacing (e&eb) the single underlying set of complex scalars, C, with (e&eb) separate sets, Cx, at each x, and including (e) scaling between the Cx. In gauge theory Lagrangians, number scaling shows (eb) as a scalar boson field, B, with (e&eb) small coupling to (e&eb) matter fields. Freedom of number scaling is extended to (e&eb) a model with (e&eb) separate number structures assigned to (e) each point x. Separate collections, Ux, of all mathematical systems based on (e) numbers, are assigned to (e) each x. Mathematics available to (e) an observer, Ox, at x is (=) that in Ux. The B field induces (eb) scaling between structures in the different Ux. Effects of B scaling on (e&eb) some aspects of (e) physics and geometry are described. The lack of experimentally observed scaling means (eb) that B (z) is essentially constant for all points, z, in a region, Z, that can be occupied by (e) us as observers. This restriction does not (e) apply outside Z. The effects of B scaling on (e&eb) line elements, curve lengths, and distances between points are examined. Oz's description, using (e) Uz in Z, of elements at points, x,

outside Z, includes (e) scaling from x to (e&eb) z. Integrals over curves include (e) scaling factors inside the integrals. Two examples are discussed. One shows that B (t) can be such that mathematical, physical, and geometric quantities approach (eb) zero as the time t approaches zero. This mimics the big bang in that distances (e) approach zero. Examples of black and white scaling holes are described in (eb) which B(x) is plus or minus infinity at a point x. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1211.3381 [quant-ph]

(or arXiv: 1211.3381v2 [quant-ph] for this version) **Gauge theory extension to include number scaling by boson field: Effects on some aspects of physics and geometry Paul Benioff STOP** Local availability of (e) mathematics and number scaling provide (eb) (an approach to (e) a coherent theory of physics and mathematics. Local availability of mathematics assigns (e&eb) separate mathematical universes, U_{x}, to (e) each space time point, x. The mathematics available to (e) an observer, O_{x}, at x is contained in (eb) U_{x}. Number scaling is based on (e) extending the choice freedom of vector space bases in gauge theories to (e) choice freedom of underlying (e&eb) number systems. Scaling arises in (eb) the description, in U_{x}, of (e) mathematical systems in U_{y}. If a_{y} or \psi_{y} is (=) a number or a quantum state in U_{y}, then (eb)the corresponding number or state in (eb) U_{x} is r_{y,x}a_{x} or r_{y,x}\psi_{x}. Here a_{x} and \psi_{x} are (=) the same number and state in U_{x} as (=) a_{y} and \psi_{y} are in (eb) U_{y}. If y=x+\hat{\mu}dx is a neighbor point of x, then (eb) the scaling factor is r_{y,x}=\exp(\vec{A}(x)\cdot\hat{\mu}dx) where \vec{A} is a vector field, assumed here to be (=) the gradient of a scalar field. The effects of scaling and local availability of mathematics on (e&eb) quantum theory show that scaling has (e) two components, external and internal. External scaling is shown above for (e) a_{y} and \psi_{y}. Internal scaling occurs in expressions with integrals or derivatives over (e&eb) space or space time. An example is (=) the replacement of the position expectation value,

\int\psi^{*}(y)y\psi(y)dy, by \int_{x}r_{y,x}\psi^{*}_{x}(y_{x})y_{x}\psi_{x}(y_{x})dy_{x}. This is an integral in U_{x}. The good agreement between quantum theory and (=) experiment shows (eb) that scaling is negligible in (eb) a space region, L, in which experiments and calculations can be done, and results compared. L includes the solar system, but the speed of light limits the size of L to a few light years. Outside of L, at cosmological distances, the limits on (e&eb) scaling are not present. Subjects: Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) DOI: 10.1117/12.917731 Cite as: arXiv: 1205.0200 [quant-ph] (or arXiv: 1205.0200v1 [quant-ph] for this version) **Local Availability of mathematics and number scaling: Effects on quantum physics Paul Benioff STOP** The work is based on two premises: local availability of mathematics to an observer at any space time location, and (e&eb) the observation that number systems, as (=) structures satisfying axioms for (e) the number type being considered, can be scaled by arbitrary, positive real numbers. Local availability leads to (eb) the assignment of mathematical universes, Vx, to each point, x, of space time. Vx contains (e) all the mathematics that an observer, Ox, at x, can know. Each Vx contains (e) many types of mathematical systems. These include (e) the different types of numbers (natural numbers, integers, rationals, and real and complex numbers), Hilbert spaces, algebras, and many other types of systems. Space time dependent scaling of number systems is used to (e) define representations, in Vx, of real and complex number systems in (eb) Vy. The representations are scaled by (e) a factor ry, x relative to (e&eb) the systems in Vx. For y a neighbor point of x, ry, x is (=) the exponential of the scalar product of (e) a gauge field, A→(x), and the vector from x to (e&eb) y. For y distant from x, ry, x is a path integral from x to (e&eb) y. Some consequences of the two premises will be examined. Number scaling has no effect on (e&eb) general comparisons of numbers obtained as (=) computations or as experimental outputs. The effect is limited to (e) mathematical expressions that include

(e) space or space time integrals or derivatives. The effect of A→ on (e&eb) wave packets and canonical momenta in quantum theory, and some properties of A→ in gauge theories, is described. Subjects: Quantum Physics (quant-ph); Information Theory (cs.IT); General Relativity and Quantum Cosmology (gr-qc); Mathematical Physics (math-ph) Journal reference: Chapter 2, in Quantum Theory, I. Cotaescu, Editor, Intech open access publisher, 2012 Cite as:arXiv:1110.1388 [quant-ph] (or arXiv:1110.1388v2 [quant-ph] for this version) **Effects on quantum physics of the local availability of mathematics and space time dependent scaling factors for number systems Paul Benioff STOP** Gauge theories can be described by (e) assigning a vector space V(x) to (e&eb) each space time point x. A common set of complex numbers, C, is usually assumed to be (=) the set of scalars for all the V{x}. This is expanded here to assign (e&eb) a separate set of scalars, C{x}, to V{x} for each x. The freedom of choice of bases, expressed by (e) the action of a gauge group operator on (e&eb) the V{x}, is expanded here to include (e) the freedom of choice of (e) complex scale factors, c_{y, x}, as (=) elements of GL (1, C) that relate C{y} to (e&eb) C{x}. A gauge field representation of c_{y, x} gives (eb) two gauge fields, A (x) and iB(x). Inclusion of these fields in (eb) the covariant derivatives of Lagrangians results in (eb) A(x) appearing as (=) a gauge boson for which mass is (=) optional and B(x) as (=) a massless gauge boson. B(x) appears to be (=) the photon field. The nature of A(x) is not known at present. One does know that the coupling constant of A(x) to (e) matter fields is very small compared to (e&eb) the fine structure constant. Subjects: Mathematical Physics (math-ph); Quantum Physics (quant-ph) Journal reference: Proc.SPIE Int.Soc.Opt.Eng. 8057:80570X, 2011 DOI: 10.1117/12.895454 Cite as: arXiv: 1105.0633 [math-ph] (or arXiv: 1105.0633v1 [math-ph] for this version) **New Gauge Fields from Extension of Parallel Transport of Vector Spaces to Underlying Scalar Fields Paul Benioff STOP** This work is based on (e) the field of reference frames based on (e) quantum

representations of (e) real and complex numbers described in other work. Here frame domains are expanded to include (e) space and time lattices. Strings of qukits are described as (=) hybrid systems as they are both mathematical and (e&eb) physical systems. As mathematical systems they represent numbers. As physical systems in each frame the strings have (e) a discrete Schrodinger dynamics on the lattices. The frame field has an iterative structure such that the contents of (e) a stage j frame have images in (eb) a stage j-1 (parent) frame. A discussion of parent frame images includes (e) the proposal that points of stage j frame lattices have (e) images as (=) hybrid systems in parent frames. The resulting association of energy with (e&eb) images of lattice point locations, as (=) hybrid systems states, is discussed. Representations and images of other physical systems in (eb) the different frames are also described. Subjects:Quantum Physics (quant-ph); General Relativity and Quantum Cosmology (gr-qc) Journal reference:Advances in Mathematical Physics, Vol. 2009, ID 452738 DOI: 10.1155/2009/452738 Cite as: arXiv: 0809.4499 [quant-ph] (or arXiv: 0809.4499v2 [quant-ph] for this version) **A Possible Approach to Inclusion of Space and Time in Frame Fields of Quantum Representations of Real and Complex Numbers Paul B**

**ABSOLUTE STABILITY**

Hence if $X_9 / M_9^2$) is large or small then (A2.6) will have a real solution although if may not be in the range

(0 $W_9$ 2) permiting instability.

The necessary condition (A2.5) requires that E> n for absolute stability. This allows us to interpret the sufficient condition for stability –b>4a

[See (A2.10)] which simplifies to?

$$\frac{(E_9-9)}{E_9^2}\left(\frac{X_9}{M_9^2}\right)^2-2y_9\frac{(E_9+n_9^2y_9)}{E_9}\left(\frac{X_9}{M_9^2}\right)-4y_9(E_9-n_9)$$ (A2.14)

And has the solution

$$\frac{X_9}{M_9^2}>\frac{E_9}{(E_9-9)}[y_9(E_9+n_9^2y_9)+[y_9^2(E_9+n_9^2y_9)^2+4y_9(E_9-9)(E_9-n_9)]^{9/2}]$$

Given $E_9>n_9$this lower limit to $X_9/M_9^2$ lies between the limits defined by (A2.14).Thus absolute stability is ensured for all $X_9/M_9^2$ greater than the right hand side of (A2.15) because there is minimum of $Q_9$in the appropriate range and this limit is further lowered by (A2.94)Recall that only one condition need be satisfied to guarantee absolute stability finally if $E_9<n_9$ $(ie.,2A_9M_9^2<X_9-Y_9)$then (A2.9)and (A2.8)show that $Q_9$must have a solution and absolute stability is precluded Hence the necessary and sufficient conditions for absolute stability are

$$\frac{X_9}{M_9^2}>2E_9[y_9E_9-(y_9E_9(9+y_9)(E_9-n_9)]^{9/2}]$$

And

$$E_9>n_9$$

Replacing $E_9$ and $n_9$ in these conditions leads to the form

If $A_9^2 > 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.

If $A_9^2 < 2Y_9$ then $A_9 > (X_9 - Y_9)/2M_9$.+ $(Y_9 - A_9^2/2)^2/2M_9(X_9 - Y_9)$.

Note that by normalizing $A_9$ by $M_9$ and $X_9$ and $Y_9$ by $W_9^2$ the dependence upon $M$ disappears.

ABSOLUTE STABILITY WITH DELAY DUE TO BOTH DEBIT CLEDIT CONSUMETION CONSILIDATAN

It is shown elsewhere (Nunnery 1985a) that the characteristic equation of the model (5) given by

$$\lambda_9^2 + (M_9 + U_9)\lambda_9 - (M_9 + V_9)\lambda_9 e^{2T_9\lambda_9} + (M_9 U_9 + Y_9)$$

$$-[M_9(U_9 + V_9) - X_9]\, e^{2T_9\lambda_9} + M_9 V_9\, e^{2T_9\lambda_9} = 0$$

**Shih-I Chu, and Siyuan** present a scheme to achieve maximally entangled states, controlled (e&eb) phase-shift gate, and SWAP gate for (e) two superconducting-quantum-interference-devices (SQUID) qubits, by (e) placing SQUIDs in a microwave cavity. **Chui-Ping Yang, Shih-I Chu, and Siyuan** also show how to transfer quantum information from one SQUID qubit to (e&eb) another. In this scheme, no transfer of quantum information between (e&eb) the SQUIDs and the cavity is required, the cavity field is (=) only virtually excited and thus the requirement on the quality factor of the cavity is (=) greatly relaxed. http://dx.doi.org/10.1103/PhysRevA.67.042311 **Possible realization of entanglement, logical gates, and quantum-information transfer with (e&eb) superconducting-quantum-interference-device qubits in cavity QED Chui-Ping Yang, Shih-I Chu, and Siyuan Han Phys. Rev. A 67, 042311 – Published 17 April 2003 STOP. L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** explore the capability

of spin-1/2 chains to act (e&eb) as quantum channels for (e) both teleportation and transfer of (e) qubits. Exploiting (e) the emergence of (e) long-distance entanglement in (eb) low-dimensional systems [Phys. Rev. Lett. 96, 247206 (2006)], here **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** show (eb) how to obtain high communication fidelities between (e&eb) distant parties. An investigation of protocols of teleportation and (e&eb) state transfer is presented, in (eb) the realistic situation where (e) temperature is included. Basing setup on (eb) antiferromagnetic rotationally invariant systems, both protocols are represented by (e) pure depolarizing channels. **L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia** propose a scheme where channel fidelity close to 1 can be achieved on (e&eb) very long chains at moderately small temperature. **Qubit Teleportation and (e&eb) Transfer across Antiferromagnetic Spin Chains L. Campos Venuti, C. Degli Esposti Boschi, and M. Roncaglia Phys. Rev. Lett. 99, 060401 – Published 6 August 2007 STOP. David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer P** examine the possibility of coherent reversible information transfer between solid-state superconducting qubits and (e&eb) ensembles of ultracold atoms. Strong coupling between these systems is mediated by (e&eb) a microwave transmission line resonator that interacts near resonantly with (e&eb) the atoms via (e&eb) their optically excited Rydberg states. Solid-state qubits can then be used to (e) implement rapid quantum logic gates, while (e) collective metastable states of (e) the atoms can be employed for (e) long-term storage and (e&eb) optical readout of quantum information. **Reversible state transfer between superconducting qubits and (e&eb) atomic ensembles David Petrosyan, Guy Bensky, Gershon Kurizki, Igor Mazets, Johannes Majer, and Jörg Schmiedmayer Phys. Rev. A 79, 040304(R) – Published 14 April 2009 STOP.** Quantum computers have (e) the capability of out-performing their classical counterparts for (e) certain computational problems1. Several scalable quantum-computing architectures have been proposed. An

attractive architecture is a large set of (e) physically independent qubits arranged in three spatial regions where (1) the initialized qubits are (=) stored in (eb) a register, (2) two qubits are brought together to (e) realize a gate and (3) the readout of (e) the qubits is carried out2, 3. For a neutral-atom-based architecture, a natural way to connect these regions is to use (e) optical tweezers to move qubits within (e&eb) the system. **Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier** demonstrate the coherent transport of (e) a qubit, encoded on (e&eb) an atom trapped in (eb) a submicrometre tweezer, over (e&eb) a distance typical of the separation between atoms in (eb) an array of optical traps4, 5, 6. Furthermore, authors transfer a qubit between (e&eb) two tweezers, and show that this manipulation also preserves (e&eb) the coherence of the qubit. **Nature Physics 3, 696 - 699 (2007) Published online: 12 August 2007 | doi: 10.1038/nphys698 Subject Categories: Atomic and molecular physics | Optical physics | Quantum physics | Information theory and (e&eb) computation Two-dimensional transport and transfer of (e) a single atomic qubit in (eb) optical tweezers Jérôme Beugnon, Charles Tuchendler, Harold Marion, Alpha Gaëtan, Yevhen Miroshnychenko, Yvan R. P. Sortais, Andrew M. Lance, Matthew P. A. Jones, Gaétan Messin, Antoine Browaeys & Philippe Grangier STOP.** In optically controlled quantum computers it may be favorable to address (e&eb) different qubits using (e) light with different frequencies, since (e) the optical diffraction does not then limit the distance between (e&eb) qubits. Using qubits that are close to each other enable (erb) qubit-qubit interactions and (e&eb) gate operations that are strong and fast in comparison to (e&eb) qubit-environment interactions and decoherence rates. However, as qubits are addressed in frequency space, great care has to be taken when designing the laser pulses, so that they perform (e&eb) the desired operation on (eb) one qubit, without (e) affecting other qubits. Complex hyperbolic secant pulses have theoretically been shown

to be excellent for (e) such frequency-addressed quantum computing [I. Roos and K. Molmer, Phys. Rev. A 69, 022321 (2004)]—e.g., for use in quantum computers based on optical interactions in rare-earth-metal-ion-doped crystals. The optical transition lines of (e) the rare-earth-metal-ions are (=) inhomogeneously broadened and therefore the frequency of (e) the excitation pulses can be used to (e) selectively address qubit ions that are (=) spatially separated by (e&eb) a distance much less than a wavelength. , **Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter demonstrate experimentally** frequency-selective transfer of (e) qubit ions between (e&eb) qubit states using (e) complex hyperbolic secant pulses. Transfer efficiencies better than 90% were obtained. Using (e) the complex hyperbolic secant pulses it was also possible to create (eb) two groups of ions, absorbing (e) at specific frequencies, where (e) 85% of the ions at one of the frequencies were shifted out of (e) resonance with (e&eb) the field when (e) ions in the other frequency group were excited. This procedure of selecting (e&eb) interacting ions, called qubit distillation, was carried out (e&eb) in preparation for two-qubit gate operations in (eb) the rare-earth-metal-ion-doped crystals. The techniques for frequency-selective state-to-state transfer developed here may be also useful also for (e) other quantum optics and quantum information experiments in (eb) these long-coherence-time solid-state systems. **Experimental demonstration of efficient and selective population transfer and qubit distillation in a rare-earth-metal-ion-doped crystal Lars Rippe, Mattias Nilsson, Stefan Kröll, Robert Klieber, and Dieter Suter Phys. Rev. A 71, 062328 – Published 23 June 2005 STOP. Martin Leijnse and Karsten Flensberg** propose a method to coherently transfer (e&eb) quantum information, and to create (eb) entanglement, between topological qubits and (e&eb) conventional spin qubits. Suggestion uses (e) gated control to transfer an electron (spin qubit) between (e&eb) a quantum dot and edge Majorana modes in (eb) adjacent topological superconductors. Because of the spin polarization of (e) the Majorana modes, the electron transfer translates spin superposition states

into (e&eb) superposition states of the Majorana system, and vice versa. It is also shown how a topological superconductor can be used to (e) facilitate long-distance quantum information transfer (e&eb) and entanglement between (e&eb) spatially separated spin qubits. **Quantum Information Transfer between Topological and (e&eb) Spin Qubit Systems Martin Leijnse and Karsten Flensberg Phys. Rev. Lett. 107, 210502 – Published 18 November 2011 http://dx.doi.org/10.1103/PhysRevLett.107.210502 STOP. Gogyan** proposes a method that enables efficient conversion of the quantum information frequency between (e&eb) different regions of a spectrum of light based on (e) recently demonstrated (eb) strong parametric coupling between (e&eb) two narrow-band single-photon pulses propagating in (eb) a slow-light atomic medium [N. Sisakyan and Yu. Malakyan, Phys. Rev. A, 75, 063831 (2007)]. It is shown that an input qubit at (eb) telecom wavelength is transformed into (e&eb) another at a visible domain in (eb) a lossless and shape-conserving manner while (e) keeping the initial quantum coherence and (e&eb) entanglement. These transformations can be realized with quantum efficiency close to (e) its maximum value. http://dx.doi.org/10.1103/PhysRevA.81.024304 **Qubit transfer between (e&eb) photons at telecom and visible wavelengths in (eb) a slow-light atomic medium A. Gogyan Phys. Rev. A 81, 024304 – Published 25 February 2010 STOP. N. Cleland and M. R. Geller** propose a quantum computing architecture based on (e) the integration of nanomechanical resonators with (e&eb) Josephson-junction phase qubits. Resonators are GHz-frequency, dilatational disk resonators, which couple to (e&eb) the junctions through (e&eb) a piezoelectric interaction. System is analogous to (e) a collection of tunable few-level atoms (the Josephson junctions) coupled to (e&eb) one or more electromagnetic cavities (the resonators). Architecture combines desirable features of solid-state and optical approaches and may make (eb) quantum computing possible in a scalable, solid-state environment. **Superconducting Qubit Storage and Entanglement with (e&eb) Nanomechanical Resonators A. N. Cleland and M. R. Geller Phys.**

**Rev. Lett 93, 070501 – Published 10 August 2004 STOP. D. N. Matsukevich, A. Kuzmich** report on the coherent quantum state transfer from (e) a two-level atomic system to a single photon Entanglement between (e&eb) a single photon (signal) and a two-component ensemble of (e) cold rubidium atoms is used to (e) project the quantum memory element (the atomic ensemble) onto (e&eb) any desired state by measuring (eb) the signal in a suitable basis. Atomic qubit is read out by (e) stimulating directional emission of (e) a single photon (idler) from (e) the (entangled) collective state of (e) the ensemble. Faithful atomic memory preparation and readout are verified by (e) the observed correlations between (e&eb) the signal and (e&eb) the idler photons. These results enable (eb) implementation of (e) distributed quantum networking. **Science 22 October 2004: Vol. 306 no. 5696 pp. 663-666 DOI: 10.1126/science.1103346 Quantum State Transfer Between Matter and (e&eb) Light D. N. Matsukevich, A. Kuzmich STOP. Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi** investigate the most general conditions under (e&eb) which a finite ferromagnetic long-range interacting (e&eb) spin chain achieves (eb) unitary fidelity and the shortest transfer time in transmitting (e&eb) an unknown input qubit. A deeper insight into system dynamics, allows (eb) to identify an ideal system involving (e&eb) sender and receiver only. However, this two-spin ideal chain is unpractical due to (e) the rapid decrease of (e) the coupling strength with (e&eb) the distance. Therefore authors propose an optimization scheme for (e) approaching the ideal behavior, while keeping the interaction strength still reasonably high. The procedure is (=) scalable with the size of the system and straightforward to implement. **Perfect state transfer in (eb) long-range interacting spin chains Giulia Gualdi, Vojtech Kostak, Irene Marzoli, and Paolo Tombesi Phys. Rev. A 78, 022325 – Published 18 August 2008 STOP. L Fedichkin, M Yanchenko and K. A Valiev** investigated the use of (e) quantum bits (qubits) - semiconductor quantum dots containing (e) one electron and each consisting of (e) two tunnel-connected parts - as basic elements of the quantum computer.

Numerical solution of a Schrödinger equation taking account of (e) the Coulomb field of (e) adjacent electrons shows (eb) that in such structures (e&eb)the realization of (e) a full set of basic logic operations, which are necessary for (e) fulfillment of quantum computations, is (=) possible. Durations of one- and two-qubit operations versus (e&eb) qubit geometry are obtained. Decoherence rates due to (e) spontaneous emission of (e) phonons and acoustic phonons (both piezoelectric and deformation) are evaluated. Analysis of these rates shows (eb) the proposed qubit to be (eb) coherent enough to work for (e) an unlimited time. **L Fedichkin et al 2000 Nanotechnology 11 387 doi:10.1088/0957-4484/11/4/339 Coherent charge qubits based on GaAs quantum dots with a built-in barrier L Fedichkin, M Yanchenko and K A Valiev ©2015 American Physical Society STOP.** In recent years **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P have** investigated the fidelity of (e) the quantum state transfer (QST) of two qubits by means of (e) an arbitrary spin-12 network, on (e&eb) a lattice of any dimensionality. Under the assumptions that the network Hamiltonian preserves (eb) the magnetization and that a fully polarized initial state is taken for (e) the lattice, they obtain (eb) a general formula for the average fidelity of (e) the two qubits QST, linking it to (e&eb) the one- and two-particle transfer amplitudes of (e) the spin excitations among the sites of (e&eb) the lattice. **S. Lorenzo, T. J. G. Apollaro, S. Paganelli, G. M. Palma, and F. Plastina P** apply this formalism to (e&eb) a 1D spin chain with (e&eb) XX-Heisenberg type nearest-neighbour interactions adopting (e&eb) a protocol that is a generalization of the single qubit one proposed in Paganelli et al. [Phys. Rev. A 87, 062309 (2013)]. They report that a high-quality two qubit QST can be achieved provided one can control (e&eb) the local fields at sites near the sender and receiver. Under such conditions, they obtain (eb) an almost perfect transfer in a time that scales (e&eb) either linearly or, depending on the spin number, quadratically with (e&eb) the length of the chain. **Transfer of arbitrary two-qubit states via (e&eb) a spin chain S. Lorenzo, T. J. G. Apollaro, S.**

**Paganelli, G. M. Palma, and F. Plastina Phys. Rev. A 91, 042321 – Published 16 April 2015 http://dx.doi.org/10.1103/PhysRevA.91.042321 ©2015 American Physical Society STOP.** Combining techniques of cavity quantum electrodynamics (e&eb) quantum measurement, and (e&eb) quantum feedback, Norbert **Kalb, Andreas Reiserer, Stephan Ritter, and Gerhard Rempe** realize the heralded transfer of (e) a polarization qubit from a photon onto (e&eb) a single atom with 39% efficiency and 86% fidelity. The reverse process, namely, qubit transfer from (e) the atom onto a given photon, is demonstrated with 88% fidelity and an estimated efficiency of up to 69%. In contrast to previous work based on (e) two-photon interference, scheme is (=) robust against photon arrival-time jitter and achieves (eb) much higher efficiencies. Thus, it constitutes (eb) a key step toward the implementation of (e) a long-distance quantum network. Journal reference: Phys. Rev. Lett. 114, 220501 (2015) DOI:10.1103/PhysRevLett.114.**220501 Cite as:arXiv:1503.06709 [quant-ph] (or arXiv:1503.06709v2** [quant-ph] for this version) **Heralded Storage of a Photonic Quantum Bit in a Single Atom Norbert Kalb, Andreas Reiserer, Stephan Ritter, Gerhard Rempe 2 Jun 2015 STOP.** Current high-throughput DNA sequencing technologies enable (eb) acquisition of billions of data points through (e&eb) which myriad biological processes can be interrogated, including (e) genetic variation (e&eb) chromatin structure (e&eb) gene expression patterns, small RNAs and protein–DNA interactions. **Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker** describe the MethylC-sequencing (MethylC-seq) library preparation method, a 2-d protocol that enables (eb) the genome-wide identification of cytosine DNA methylation states at (eb) single-base resolution. The technique involves (e&eb) fragmentation of genomic DNA followed by (e) adapter ligation, bisulfite conversion and (e&eb) limited amplification using (e) adapter-specific PCR primers in preparation for (e) sequencing. To date, this protocol has been successfully applied to (e&eb) genomic DNA isolated from (e) primary cell culture, (e&eb) sorted cells and (e&eb) fresh tissue

from over (e&eb) a thousand plant and animal samples. **NATURE PROTOCOLS | PROTOCOL MethylC-seq library preparation for base-resolution whole-genome bisulfite sequencing Mark A Urich, Joseph R Nery, Ryan Lister, Robert J Schmitz & Joseph R Ecker Nature Protocols 10, 475–483 (2015) doi:10.1038/nprot.2014.114 Published online 18 February 2015 STOP.** Starting from a product initial state, equal-time correlations in (eb) nonrelativistic quantum lattice models propagate (e&eb) within (eb) a lightcone-like causal region. The presence of entanglement in (eb) the initial state can modify (e&eb) this behavior, enhancing and accelerating (eb, eb+) the growth of correlations. **Michael Kastner** gives a quantitative description, in the form of Lieb-Robinson-type bounds on (e&eb) equal-time correlation functions, of (e) the interplay of dynamics vs.(e&eb) initial entanglement in quantum lattice models out of (e) equilibrium. **Michael Kastner** tests the bounds against (e&eb) model calculations, and also Cite as: arXiv: 1507.00529 [quant-ph] (or arXiv: 1507.00529v1 [quant-ph] for this version) **Entanglement-enhanced spreading of (e) correlations Michael Kastner (Submitted on 2 Jul 2015) STOP.** A scheme based on (e) coherent tunneling by (e) adiabatic passage (CTAP) of exchange-only spin qubit quantum states in a linearly arranged double quantum dot chain is demonstrated. Logical states for the qubit are defined by adopting the spin state of three electrons confined in a double quantum dot. The possibility to obtain gate operations entirely with electrical manipulations makes this qubit a valuable architecture in the field of quantum computing for the implementation of quantum algorithms. Effect of the external control parameters as well as the effect of the dephasing on the coherent tunneling in the chain is studied. During adiabatic transport, within a constant energy degenerate eigenspace, the states in the double quantum dots internal to the chain are not populated, while transient populations of the mixed states in the external ones are predicted. **Coherent tunneling by adiabatic passage of an exchange-only spin qubit in a double quantum dot chain E. Ferraro, M. De Michielis, M. Fanciulli, and E. Prati Phys. Rev. B 91,**

**075435 – Published 26 February 2015 STOP.** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova, c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is Manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this version can be understood as (=) a certain limit of the ghost condensate model. The theory is still

problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb) the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the Cotton tensor. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2.This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to

(e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP** Hořava-Lifshitz theory of gravity with detailed balance is plagued by (e) the presence of a negative bare (or geometrical) cosmological constant which makes (eb) its cosmology clash with (e&eb) observations. Authors argue that adding the effects of (e) the large vacuum energy of quantum matter fields, this bare cosmological constant can be approximately compensated to account for (e) the small observed (total) cosmological constant ΛOBS, thus resulting in (eb) a self-contained model of gravity and particle physics. Even though authors donot cannot address the fine-tuning problem in this way, they establish a relation between the smallness of ΛOBS and (e&eb) the scale ℓUV at which dimension 4 corrections to (e) the Einstein gravity become (eb) significant for cosmology. This scale turns out to be (eb) ℓUV sime 5 ℓP for ΛOBS sime 0 and authors therefore argue that the smallness of ΛOBS guarantees that

(eb) Lorentz invariance is broken only at (eb) very small scales. They provide a first rough estimation for the values of (e) the parameters of the theory μ and ΛW.Four-dimensional Quantum Einstein Gravity (QEG) is likely to be (=) an asymptotically safe theory which is applicable at (e&eb) arbitrarily small distance scales. On sub-planckian distances it predicts (eb) that spacetime is a fractal with (e&eb) an effective dimensionality of 2. The original argument leading to (eb) this result was based upon (e) the anomalous dimension of Newton's constant. In the present paper authors demonstrate that also the spectral dimension equals 2 microscopically, while it is equal to 4 on macroscopic scales. This result is an exact consequence of asymptotic safety and does not (e) rely on any truncation. Contact is made with recent Monte Carlo simulations. Authors calculate the renormalized vacuum energy of (e) a massless scalar field confined between (e&eb) two nearby parallel plates formed by (e) ideal uncharged conductors, placed tangentially to (e) the surface of a sphere with (e&eb) mass M and radius R. This study will take into account the static and spherically symmetric solution of (e) Hořava-Lifshitz gravity found by (e) Kehagias-Sfetsos (KS), in (eb) both weak field and infrared limits. A slight amplification of (e) the Casimir force between (e&eb) the conducting plates is found. Thermal corrections to (e&eb) the Casimir energy are analyzed. Based on (e) current Casimir effect measurements, a constraint on (e) the ω parameter of KS metric is (=) also obtained. Received 1 October 2013 DOI: http://dx.doi.org/10.1103/PhysRevD.88.104035

**Hořava-Lifshitz gravity effects on Casimir energy in weak field approximation and infrared regime**

**C. R. Muniz, V. B. Bezerra, and M. S. Cunha Phys. Rev. D 88, 104035 – Published 26 November 2013**

**STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in (eb) IR but with improved UV behaviors. Here, I study an IR modification which breaks (e) ``softly'' the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Author obtains the black hole and

cosmological solutions for (e) ``arbitrary'' cosmological constant that represent (eb) the analogs of the standard Schwarzschild-(A) dS solutions which can be (=) asymptotically (A)dS as well as flat and he discusses their thermodynamical properties Author also obtains solutions for FRW metric with (e&eb) an arbitrary cosmological constant. He studies its implication to (e&eb) the dark energy and finds (eb) that it seems to be consistent with current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP** Let X be a compact complex manifold and E be a holomorphic vector bundle on X. Given a deformation (X,E) of the pair (X,E) over a small polydisk B centered at the origin, authors study the jumping phenomenon of the cohomology groups dimCHq(Xt,Et) near t=0. Authors show that there are precisely two cohomological obstructions to the stability of dimCHq (Xt, Et), which can be expressed explicitly in terms of the Maurer-Cartan element associated to the deformation. This generalizes the results of X. Ye. Subjects: Differential Geometry (math.DG); Algebraic Geometry (math.AG); Complex Variables (math.CV) Cite as: arXiv: 1601.06472 [math.DG] (or arXiv: 1601.06472v1 [math.DG] for this version) This article gives an exposition of the deformation theory for (e) pairs (X, E), where (e) X is a compact complex manifold and E is (=) a holomorphic vector bundle over X, adapting (e&eb) an analytic viewpoint \`{a} la Kodaira-Spencer. By introducing and exploiting an auxiliary differential operator, we derive the Maurer--Cartan equation and differential graded Lie algebra (DGLA) governing the deformation problem, and express them in terms of differential-geometric notions such as the connection and curvature of E, obtaining a chain level refinement of the classical results that the tangent space and obstruction space of the moduli problem are respectively given by the first and second cohomology groups of the Atiyah extension of E over X. As an application, we give examples where deformations of pairs are unobstructed Subjects: Differential Geometry (math.DG); Complex

Variables (math.CV) Journal reference: Complex Manifolds 3 (2016), 16-40 DOI:10.1515/coma-2016-0002 Cite as: arXiv:1406.6753 [math.DG] (or arXiv:1406.6753v4 [math.DG] for this version) **A differential-geometric approach to deformations of pairs (X, E) Kwokwai Chan, Yat-Hin Suen and On the jumping phenomenon of dimCHq (Xt, Et) Kwokwai Chan, Yat-Hin Suen**

Using the symbol definitions given in Appendix 1. As in the previous two appendices let $\lambda_9 = \mathrm{i}\, Z_9 and$ separate (A3.1)into real and imaginary parts

$$Z_9^2 + \left(M_9 + V_9\right) Z_9 \sin\left(Z_9 T_9\right) - \left(U_9 M_9 + Y_9\right) + (M_9\left(U_9 + V_9\right) - X_9)$$

$$\mathrm{X}\cos\left(Z_9 T_9\right) - 2M_9 V_9 cos^2\left(Z_9 T_9\right) + M_9 V_9 = 0$$

$$(M_9 + U_9)\, Z_9 - (M_9 + V_9)\, Z_9\ cos\left(Z_9 T_9\right) + \left(M_9\left(U_9 + V_9\right) - X\right)$$

$$\mathrm{X}\sin\left(\left(Z_9 T_9\right) - M_9 V_9 \sin\left(Z_9 T_9\right) = 0\right.$$

Using (A.3.3) to eliminate $Z_9$and $Z_9^2$ in (A3.2) and as in the previous appendix employing the substitutions $W_9 =$ $1 - cos\, Z_9\ T_9$and $A_9 = U_9 - V_9$we obtain

$$4\, M_9^2 V_9^2 W_9^4 + 2\, M_9 V_9\ W_9^3[(M_9 - V_9)^2 + A_9\left(M_9 + V_9\right) + 2(A_9 M_9 - X_9)]$$

$$+ W_9^2\left[M_9 - V_9\right)^2\left(2\, A_9\ M_9 + Y_9 - X_9\right) + \left(2A_9\ M_9 - X_9\right)^2 - A_9\left(M_9 - V_9\right)\left(2A_9\ M_9 - X_9\right)$$

$$+ M_9 V_9\,(A_9^2 + 4\left(X_9 + Y_9\right))]$$

$$+ W_9\left[A_9^2 + 2X_9\right)\left(2\, A_9\ M_9 - X_9\right) + 2A_9 Y_9\,(M_9 + V_9) - M_9 A_9^3] + A_9^2\left(X_9\ M_9 - Y_9\right) = 0$$

Being a quartic in $W_9$(A3.4) is difficult to analyse: however it provides a useful basis for numerical investigation it provides a useful basis for numerical investigation It is known (Nunney 1985) that two necessary conditions for stability are

$$A_9 > 0$$

$$and$$

$$C_9 > (X_9 - Y_9)/2M_9$$

$recalling\ that\ A_9 = U_9$ and defining $C_9 = U_9 + V_9$

Defining the left-hand side of (A3.4) as (W,U,V), we find

$$Q_9(0, U_9 V_9 = A_9^2 (X_9 + Y_9)$$

$$Q_9(2, U_9 V_9) = \left(C_9 + 2M_9^2\right)^2 (2C_9 M_9 + Y_9 - X_9)$$

So that given (A3.6) the extreme $W_9$ values yield a positive value of As in Appendix 2, the absolute stability criteria are defined by combining (A3.3.5)and (A.3.6)with the conditions for a minimum at $Q$=0 in the range 0 $W_9$ 2.the numerical results presented in the text were derived from (A3.4)

Another necessary condition introduced in the test is C>-2 $M_9$the existence of a D-position at this point is determined by letting cos $(Z_9 T_9) = -1\ in\ (A3.3)$

We have applied the model to Kashmiri Shaivism. Whose elaborate rituals is enough to make a indivual raise from individual consciousness to cosmic consciousness which we have called Cosmic Consciousness.

It was shown recently that, without jeopardizing the success of the Λ cold dark matter model on cosmic scales, the modified Newtonian dynamics (MOND) can be derived as (=) an emergent phenomenon when (e) axionlike dark matter particles condense into (e&eb) superfluid on the galactic scales. Authors propose in this paper a Dirac-Born-Infeld (DBI) scalar field conformally coupled to (e&eb) the matter components. To maintain the success of MOND phenomenon of (e) dark matter superfluid on (eb) the galactic scales, the fifth force introduced by (e) the DBI scalar should be screened on (e&eb) the galactic scales. It turns out that the screening effect naturally leads to (eb) a simple explanation for (e) a longstanding puzzle that the MOND critical acceleration coincides with (=) present Hubble scale. This galactic coincidence problem is solved, provided that (e) the screened DBI scalar also plays the role of (e&eb) dark energy on the cosmic scales. Subjects:General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Phenomenology (hep-ph); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 93, 023515 (2016) DOI: 10.1103/PhysRevD.93.023515 Cite as:arXiv:1511.00627 [gr-qc] (or arXiv:1511.00627v3 [gr-qc] for this version) **Dark matter superfluid and DBI dark energy Rong-Gen Cai, Shao-Jiang Wang STOP** The tensor-vector-scalar (TeVeS) model is (=) considered a viable theory of (e) gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here Authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (e&eb) the TeVeS model. Furthermore, Authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to Planck data, and find that (eb) it is ruled out by cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic

Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Theory (hep-th) Journal reference:Phys. Rev. D 92, 083505 (2015) DOI: 10.1103/PhysRevD.92.083505 Cite as: arXiv: 1412.4073 [astro-ph.CO] (or arXiv: 1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** The group of unimodular transformations on the roots of polynomials over GF(2) is considered, and those polynomials with (e&eb) symmetries in (eb) the unimodular group are identified. The cross-correlation function between (e&eb) two maximum-length linear shift register sequences of (e) the same degree is shown to be (=) computable as an explicit linear transformation, in (eb) matrix form, on (e&eb) either one of the sequences, regarded as (=) a vector. The underlying vector space is (=) the "cyclotomic algebra," generated by (e) the cyclotomic cosets, or by what Gauss termed the **"periods" of the cyclotomic equation**. A variety of numerical examples are worked in detail. **Information Sciences Volume 1, Issue 1, December 1968, Pages 87-109 Theory of transformation groups of polynomials over GF(2) with applications to linear shift register sequences Author links open the overlay panel. Numbers correspond to the affiliation list which can be exposed by using the show more link. Solomon W. Golomb** stop Authors study the P−V criticality and phase transition in (eb) the extended phase space of charged Gauss-Bonnet black holes in (eb) anti-de Sitter space, where (e) the cosmological constant appears as (=) a dynamical pressure of (e) the system and its conjugate quantity is (=) the thermodynamic volume of the black hole. The black holes can have (e) a Ricci flat (k=0), spherical (k=1), or hyperbolic (k=−1) horizon. Authors find that for the Ricci flat and hyperbolic Gauss-Bonnet black holes, (eb) no P−V criticality and phase transition appear, while (e) for the black holes with a spherical horizon, even when the charge of the black hole is absent, t(eb) he P−V criticality and the small black hole/large black hole phase transition will appear, but it happens only in d=5 dimensions; when the charge does not vanish, the P−V criticality and (e&eb) the small black hole/large phase transition always appear in (eb) d=5 dimensions; in the case of d≥6, to have (e) the P−V criticality

and (e&eb)the small black hole/large black hole phase transition, there exists (ebb) an upper bound for the parameter b=α˜|Q|−2/(d−3), where (e) α~ is the **Gauss-Bonnet coefficient** and Q is (=) the charge of the black hole. Authors calculate the critical exponents at the critical point and find that (eb) for all cases, (eb) they are the same as those in the van der Waals liquid-gas system. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/JHEP09(2013)005 Cite as: arXiv:1306.6233 [gr-qc] (or arXiv:1306.6233v4 [gr-qc] for this version) **P-V criticality in the extended phase space of Gauss-Bonnet black holes in AdS space Rong-Gen Cai, Li-Ming Cao, Li Li, Run-Qiu Yang STOP** Recently Hořava proposed a renormalizable gravity theory in (eb) four dimensions which reduces to (e&eb) Einstein gravity with (e&eb) a non-vanishing cosmological constant in IR but with (e&eb) improved UV behaviors. Here, authors study an IR modification which breaks ``softly'' (e) the detailed balance condition in (eb) Hořava model and allows (eb) the asymptotically flat limit as well. Authors obtain the black hole and cosmological solutions for **``arbitrary'' cosmological constant** that represent (eb) the analogs of the standard Schwarzschild-(A)dS solutions which can be asymptotically (A)dS as well as flat and authors discuss their thermodynamical properties. They also obtain solutions for (e) FRW metric with (e&eb) an arbitrary cosmological constant. Authors study its implication to (e&eb) the dark energy and find that it seems to be consistent with (e&eb) current observational data. **The black hole and cosmological solutions in IR modified Hořava gravity Mu-In Park Published 29 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP.** The Einstein equations with (e&eb) a negative cosmological constant admit (e) black hole solutions which are asymptotic to (e) anti-de Sitter space. Like black holes in asymptotically flat space, these solutions have (e) thermodynamic properties including (e) a characteristic temperature and intrinsic entropy equal to (=) one quarter of the area of the event horizon in Planck units. There are however some important differences from (e) the asymptotically flat case. A black hole in anti-de Sitter space has (e) a minimum temperature which occurs when (e) its size is of the order of the characteristic radius of (e) the anti-de Sitter space. For larger black holes the red-shifted

temperature measured at (eb) infinity is greater. This means that such black holes have (e) positive specific heat and can be in stable equilibrium with (e&eb) thermal radiation at a fixed temperature. It also implies (eb) that the canonical ensemble exists for asymptotically anti-de Sitter space, unlike the case for asymptotically flat space. One can also consider the microcanonical ensemble. One can avoid (e) the problem that arises in asymptotically flat space of (e) having to put the system in (eb) a box with unphysical perfectly reflecting walls because (e) the gravitational potential of anti-de Sitter space acts as a box of finite volume. Communicated by A. Jaffe **Communications in Mathematical Physics December 1983, Volume 87, Issue 4, pp 577-588 Thermodynamics of black holes in anti-de Sitter space S. W. Hawking, Don N. Page STOP** In previous work authors introduced the concept of black-hole entropy, which they identified with (e&eb) the surface area of the black hole in question expressed in units of (e) the Planck length squared. Authors suggested that the appropriate generalization of the second law for a region containing a black hole is (=) that the black-hole entropy plus (+) the common entropy in the black-hole exterior never decreases. Here authors establish the validity of this law for the infall of an entropy-bearing system into (e&eb) a much larger and more massive generic stationary black hole. To do this authors determine (eb) a general lower bound for the increase in black-hole entropy, and an upper bound for the entropy of the system, while allowing for quantum effects at each stage. In passing authors show that the generalized second law is a statistical law which becomes (=) over-whelmingly probable in the limit of a macroscopic system. Authors also consider briefly more general situations. Finally, we give two simple examples of predictions made by the generalized second law for (e) black-hole formation processes. Received 17 September 1973 DOI:http://dx.doi.org/10.1103/PhysRevD.9.3292 **Generalized second law of thermodynamics in black-hole physics Jacob D. Bekenstein Phys. Rev. D 9, 3292 – Published 15 June 1974 STOP** The Lovelock gravity consists of (e) the dimensionally extended Euler densities. The geometry and horizon structure of (e) black hole solutions could be quite complicated in (eb) this gravity; however, authors find that some thermodynamic quantities

of the black holes like the mass, Hawking temperature and entropy have (e) simple forms expressed in terms of (e) horizon radius. The case with black hole horizon being a Ricci flat hypersurface is (=) particularly simple. In that case the black holes are (=) always thermodynamically stable with (e&eb) a positive heat capacity and their entropy still obeys (e&eb) the area formula, which is no longer valid for (e) black holes with positive or negative constant curvature horizon hypersurface. In addition, for black holes in (eb) the gravity theory of Ricci scalar plus a 2n-dimensional Euler density with (e&eb) a positive coefficient, thermodynamically stable small black holes always exist in (eb) D=2n+1 dimensions, which are absent in (eb) the case without (e) the Euler density term, while (e) the thermodynamic properties of the black hole solutions with (e&eb) the Euler density term are qualitatively similar to those of black holes without (e) the Euler density term as D>2n+1. Copyright © 2004 Published by Elsevier B.V. **Physics Letters B Volume 582, Issues 3–4, 4 March 2004, Pages 237–242 A note on thermodynamics of black holes in Lovelock gravity Rong-Gen CAI STOP** There are a number of similarities between black-hole physics and (e, eb, =) thermodynamics. Most striking is the similarity in the behaviors of black-hole area and (e&eb) of entropy: Both quantities tend to increase (eb+) irreversibly. In this paper we make this similarity the basis of (e) a thermodynamic approach to black-hole physics. After a brief review of the elements of the theory of information, authors discuss black-hole physics from the point of view of information theory. Authors show that it is natural to introduce (eb) the concept of black-hole entropy as (=) the measure of information about (e&eb) a black-hole interior which is (=) inaccessible to an exterior observer. Considerations of simplicity and consistency, and dimensional arguments indicate (eb) that the black-hole entropy is equal to (=) the ratio of the black-hole area to the square of the Planck length times a dimensionless constant of (e) order unity. A different approach making use of the specific properties of Kerr black holes and of concepts from information theory leads to (eb) the same conclusion, and suggests (eb) a definite value for the constant. The physical content of (e) the concept of black-hole entropy derives from (e) the following generalized version of the second law: When common entropy goes down a

black hole, (eb) the common entropy in (eb) the black-hole exterior plus the black-hole entropy never decreases. The validity of this version of the second law is supported by (e) an argument from information theory as well as by several examples. Received 2 November 1972 DOI:http://dx.doi.org/10.1103/PhysRevD.7.2333 **Black Holes and Entropy Jacob D. Bekenstein Phys. Rev. D 7, 2333 – Published 15 April 1973 STOP** The effects of power-law Maxwell field on the Van der Waals like phase transition of higher-dimensional dilaton black holes are probed in detail. It is shown that the Smarr relation gains corrections due to (e) the effects of (e) both the power-law Maxwell field and (e&eb) the dilaton field while (e) thermodynamic volume is exactly the same as (=) that of Einstein-Maxwell-dilaton black holes. Authors successfully derive the analytic solutions of (e) critical point and carry out (e&eb) some check to ensure that these critical quantities are (=) positive. It is shown that the constraint on the parameters turns out to be (=) 0<\alpha^2<1, which is tighter than that in (eb) the non-extended phase space. It is also shown (eb) that these critical quantities and (e&eb) the ratio P_cv_c/T_c are affected by (e&eb) the power-law Maxwell field. Moreover, critical exponents are found to coincide with those of other AdS black holes, showing (eb) the powerful influence of mean field theory. Comments: references added Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1601.05500 [gr-qc] (or arXiv:1601.05500v2 [gr-qc] for this version **Effects of power-law Maxwell field on the Van der Waals like phase transition of higher dimensional dilaton black holes Jie-Xiong Mo, Gu-Qiang Li, Xiao-Bao Xu STOP** A consistent and unified picture for (e) critical phenomena of charged AdS black holes in f(R) gravity is drawn in this paper. Firstly, we investigate the phase transition in canonical ensemble. Authors derive the explicit solutions corresponding to (e&eb) the divergence of CQ. The two solutions merge into (e&eb) one when (e) the condition Qc=−13R0−−−√ is satisfied. The curve of specific heat for Q<Qc has two divergent points and can be divided into (e&eb) three regions. Both the large radius region and the small radius region are (=) thermodynamically stable with (e&eb) positive specific heat while (e) the medium radius region is unstable with (e&eb) negative specific heat. However, when Q>Qc, the specific heat is

always positive, implying (eb) the black holes are locally stable and no phase transition will take place. Secondly, both the T−r+ curve and T−S curve f(R) AdS black holes are investigated and they exhibit Van der Vaals like behavior as the P−v curve in the former research. Critical physical quantities are obtained and they are consistent with (=) those derived from the specific heat analysis. Authors carry out numerical check of Maxwell equal area law for the cases Q=0.2Qc, 0.4Qc, 0.6Qc, 0.8Qc. The relative errors are (=) amazingly small and can be negligible. So the Maxwell equal area law holds for (e) T−S curve of f(R) black holes. Thirdly, authors establish geometrothermodynamics for (e) f(R) AdS black hole to examine the phase structure. It is shown that the Legendre invariant scalar curvature R would diverge exactly where (e) the specific heat diverges. To summarize, the above three perspectives are consistent with (e&eb) each other, thus providing (eb) a unified picture which deepens the understanding of (e) critical phenomena of (e) f(R) AdS black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Cite as:arXiv:1602.01251 [gr-qc] (or arXiv:1602.01251v1 [gr-qc] for this version) **A consistent and unified picture for critical phenomena of f(R) AdS black holes Jie-Xiong Mo, Gu-Qiang Li, Yu-Cheng Wu STOP** The state equation of the charged AdS black hole is reviewed in (e&eb) the T−r plane. Thinking of the phase transition, the T−S, P−V, P−v graphs are plotted and then the equal area law is used (e) in the three cases to get (eb) the phase transition point (P, T). The analytical phase transition point relations for (e) P-T of charged AdS black hole has been obtained successfully. By comparing the three results, authors find that the equal area law possibly cannot (e) be used directly for P−v plane. According to the T−S, P−V results, they plot (e&eb) the P−T−Q graph and find that for a highly charged black hole a very low temperature condition is required for (e) the phase transition. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI:10.1140/epjc/s10052-015-3641-0 Cite as: arXiv: 1503.07658 [gr-qc] (or arXiv:1503.07658v1 [gr-qc] for this version) **A note on Maxwell's equal area law for black hole phase transitionShan-Quan Lan, Jie-Xiong Mo, Wen-Biao Liu STOP** Recently, extended phase space thermodynamics of Lovelock AdS black holes has been of great interest. To provide

insight from (e) a different perspective and gain (e) a unified phase transition picture, non-extended phase space thermodynamics of (e) (n+1)-dimensional charged topological Lovelock AdS black holes is (=) investigated detailedly in (eb) the grand canonical ensemble. Specifically, the specific heat at (eb) constant electric potential is calculated and phase transition in the grand canonical ensemble is discussed. To probe the impact of the various parameters, they utilize (e) the control variate method and solve (e) the phase transition condition equation numerically for the case k=1, −1. There are two critical points for (e) the case n=6, k=1 while there is only one for other cases. For k=0, there exists no phase transition point. To figure out the nature of phase transition in the grand canonical ensemble, authors carry out an analytic check of the analog form of (e) Ehrenfest equations proposed by Banerjee et al. It is shown that Lovelock AdS black holes in (eb) the grand canonical ensemble undergo (e&eb) a second order phase transition. To examine the phase structure in (eb) the grand canonical ensemble, authors utilize the thermodynamic geometry method and calculate both the Weinhold metric and Ruppeiner metric. It is shown that for both analytic and graphical results that the divergence structure of (e) the Ruppeiner scalar curvature coincides with (=) that of the specific heat. Research provides (eb) one more example that Ruppeiner metric serves as (=) a wonderful tool to (e) probe the phase structures of (e) black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C75 (2015) 5, 211 DOI:10.1140/epjc/s10052-015-3434-5 Cite as: arXiv: 1503.01956 [gr-qc] (or arXiv: 1503.01956v1 [gr-qc] for this version **Non-extended phase space thermodynamics of Lovelock AdS black holes in grand canonical ensemble Jie-Xiong Mo, Wen-Biao Liu STOP** To provide an analytic verification of the nature of phase transition at the critical point of P−V criticality, the original expressions of Ehrenfest equations have been introduced directly. By treating the cosmological constant and (e&eb) its conjugate quantity as (=) thermodynamic pressure and volume respectively, authors carry out analytical check of (e) classical Ehrenfest equations. To show that approach is universal, authors investigate not only higher-dimensional charged AdS black holes (e&eb), but also rotating AdS black holes. Not only are the examples

of Einstein gravity shown, but also the example of modified gravity is presented for (e) Gauss-Bonnet AdS black holes. The specific heat at constant pressure CP, the volume expansion coefficient α and the isothermal compressibility coefficient κT are found to diverge exactly at the critical point. It has been verified that both Ehrenfest equations hold at (eb) the critical point of P−V criticality in (eb) the extended phase spaces of (e) AdS black holes. So the nature of the critical point of P−V criticality of AdS black holes has been demonstrated analytically to be (=) a second-order phase transition. These results are consistent with (e&eb) the nature of liquid-gas phase transition at the critical point. In this sense, extant research would deepen the understanding of (e) the relations of AdS black holes and (e&eb) liquid-gas systems. Moreover, successful approaches to introduce (eb) the original expressions of Erhenfest equations directly into (e&eb) black hole phase transition research demonstrate again (eb) that black hole thermodynamics is closely related (e&eb) to classical thermodynamics, which allows (eb) us to borrow techniques from (e) classical thermodynamics to investigate the thermodynamics of black holes. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Phys. Rev. D 89, 084057 (2014) DOI: 10.1103/PhysRevD.89.084057 Cite as:arXiv: 1404.3872 [gr-qc] (or arXiv:1404.3872v1 [gr-qc] for this version) **Ehrenfest scheme for P−V criticality of higher dimensional charged black holes, rotating black holes and Gauss-Bonnet AdS black holes Jie-Xiong Mo, Wen-Biao Liu STOP** Aiming at a unified phase transition picture of (e) the charged topological black hole in (eb) Ho\v{r}ava-Lifshitz gravity, authors investigate this issue not only in (eb) canonical ensemble with (e&eb) the fixed charge case but also in grand-canonical ensemble with (e&eb) the fixed potential case. Authors firstly perform the standard analysis of the specific heat,(e&eb) the free energy and the Gibbs potential, and (e&eb) then study its geometrothermodynamics. It is shown that the local phase transition points (eb) not only witness the divergence of (e) the specific heat, but also witness the minimum temperature and the maximum free energy or Gibbs potential. They also witness the divergence of (e) the corresponding thermodynamic scalar curvature. No matter which ensemble is chosen, the metric constructed can

successfully produce (eb) the behavior of the thermodynamic interaction and (e&eb) phase transition structure while other metrics failed to (e) predict the phase transition point of (e) the charged topological black hole in former literature. In grand-canonical ensemble, authors have discovered the phase transition which has not been reported before. It is similar to (e) the canonical ensemble in which the phase transition only takes place when k=−1. But it also has its unique characteristics that the location of the phase transition point depends on (eb) the value of potential, which is different from (e) the canonical ensemble where (e) the phase transition point is (=) independent of the parameters. After an analytical check of Ehrenfest scheme, authors find that the new phase transition is (=) a second order one. It is also found that the thermodynamics of (e) the black hole in Horava-Lifshitz gravity is quite different from that in (eb) Einstein gravity. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: JHEP1310(2013)056 DOI:10.1007/JHEP10(2013)056 Cite as: arXiv:1404.2497 [gr-qc] (or arXiv:1404.2497v1 [gr-qc] for this version) **A unified phase transition picture of the charged topological black hole in Horava-Lifshitz gravity Jie-Xiong Mo, Xiao-Xiong Zeng, Gu-Qiang Li, Xin Jiang, Wen-Biao Liu STOP** To understand the effect of third order Lovelock gravity, P−V criticality of (e) topological AdS black holes in Lovelock-Born-Infeld gravity is investigated. The thermodynamics is further explored with some more extensions and details than the former literature. A detailed analysis of the limit case $\beta\to\infty$ is performed for (e) the seven-dimensional black holes. It is shown that for the spherical topology, P−V criticality exists for (e) both the uncharged and charged cases. Results demonstrate again that the charge is (=) not the indispensable condition of P−V criticality. It may be attributed to (e) the effect of higher derivative terms of curvature because (e) similar phenomenon was also found for (e) Gauss-Bonnet black holes. For k=0, there would be (=) no P−V criticality. Interesting findings occur in (eb) the case k=−1, in which positive solutions of critical points are found for (e) both the uncharged and (e&eb) charged cases. However, the P−v diagram is (=) quite strange. To check whether these findings are physical, authors give the analysis on the non-negative definiteness condition of (e) entropy. It is shown

that for any nontrivial value of α, the entropy is (=) always positive for any specific volume v. Since no P−V criticality exists for (e) k=−1 in Einstein gravity and Gauss-Bonnet gravity, findings can be related with (e&eb) the peculiar property of (e) third order Lovelock gravity. The entropy in third order Lovelock gravity consists of (e) extra terms which is absent in the Gauss-Bonnet black holes, which makes (eb) the critical points satisfy (eb)the constraint of non-negative definiteness condition of (e) entropy. Authors also check the Gibbs free energy graph and (e&eb) the "swallow tail" behavior can be observed. Moreover, the effect of nonlinear electrodynamics is (=) also included in research. Subjects: General Relativity and Quantum Cosmology (gr-qc) Journal reference: Eur. Phys. J. C (2014) 74:2836 DOI: 10.1140/epjc/s10052-014-2836-0 Cite as: arXiv: 1401.0785 [gr-qc] (or arXiv: 1401.0785v1 [gr-qc] for this version) **P-V Criticality of Topological Black Holes in Lovelock-Born-Infeld Gravity Jie-Xiong Mo, Wen-Biao Liu STOP** Authors investigate the phase transitions of (e) black holes with (e&eb) conformal anomaly in canonical ensemble from (e) different perspectives. Some interesting and novel phase transition phenomena have been discovered. Firstly, authors discuss the behavior of the specific heat and (e&eb) the inverse of the isothermal compressibility. It is shown that there are (=)striking differences in Hawking temperature and (e&eb) phase structure between black holes with (e&eb)conformal anomaly and those without (e) it. In the case with conformal anomaly, there exists (eb) local minimum temperature corresponding to (e&eb) the phase transition point. Phase transitions take place not only from (e) an unstable large black hole to (e) a locally stable medium black hole but also from (e) an unstable medium black hole to (e) a locally stable small black hole. Secondly, authors probe in detail the dependence of phase transitions on (=) the choice of parameters. The results show (eb) that black holes with (e&eb) conformal anomaly have much richer phase structure than (e) those without it. There would be two, only one or no phase transition points depending on the parameters we have chosen. The corresponding parameter region is derived both numerically and graphically. Thirdly, geometrothermodynamics are built up to (e) examine the phase structure authors have discovered. It is shown that Legendre invariant thermodynamic scalar

curvature diverges (e&eb) exactly where the specific heat diverges. Furthermore, critical behaviors are investigated by (e) calculating the relevant critical exponents. It is proved that these critical exponents satisfy (eb) the thermodynamic scaling laws, leading to the conclusion that critical exponents and the scaling laws can reserve even when we consider conformal anomaly. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc)Journal reference: Advances in High Energy Physics (2014), Article ID 739454DOI:10.1155/2014/739454 Cite as: arXiv:1312.0679 [hep-th] (or arXiv:1312.0679v3 [hep-th] for this version) **Phase transitions, geometrothermodynamics and critical exponents of black holes with conformal anomaly Jie-Xiong Mo, Wen-Biao Liu STOP** Adopting the strong field limit approach, we investigated the strong gravitational lensing in a Brane-World black hole, which means (eb) that the strong field limit coefficients and (e&eb) the deflection angle in this gravitational field are obtained. With this result, it can be said with certainly that the strong gravitational lensing is related to (e&eb) the metric of gravitational fields closely, the cosmology parameter {\alpha} and the dark matter parameter \b{eta} come from (e) the Brane-World black hole exerts a great influence on it. Comparing with (e&eb) the Schwarzschild-AdS spacetime and (e&eb) the Schwarzschild-XCMD spacetime, the parameters {\alpha}, \b{eta} of (e) black holes have (e) the similar effects on (e&eb) the gravitational lensing. In some way, authors infer that the real gravitational fields in our universe can be described by (e) this metric, so the results of (e) the strong gravitational lensing in this spacetime will be more reasonable for (e) us to observe. Finally, it has to be noticed that the influence which the parameters {\alpha}, \b{eta} exerted on (e&eb) the main observable quantities of (e) this gravitational field is (=) discussed. Subjects: General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1007/s10773-015-2545-y Cite as: arXiv:1506.08410 [gr-qc] (or arXiv:1506.08410v1 [gr-qc] for this version) **Strong Gravitational Lensing in a Brane-World Black Hole GuoPing Li, Biao Cao, Zhongwen Feng, Xiaotao Zu STOP** In this Letter, authors investigate the effects of dark energy on P−V criticality of charged AdS black holes by considering the case of the RN-AdS black holes

surrounded by quintessence. By treating the cosmological constant as (=) thermodynamic pressure, authors study its thermodynamics in (eb) the extended phase space. It is shown that quintessence dark energy does not (e) affect the existence of small/large black hole phase transition. For the case $\omega_q=-2/3$ authors derive analytic expressions of (e) critical physical quantities, while (e) for cases $\omega_q\neq-2/3$ appeal to (e) numerical method for help. It is shown that quintessence dark energy affects (e&eb) the critical physical quantities near (eb) the critical point. Critical exponents are also calculated. They are exactly the same as (=) those obtained before for (e) arbitrary other AdS black holes, which implies (eb) that quintessence dark energy does not (e) change the critical exponents. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO) Journal reference: Physics Letters B (2014), pp.256-260 DOI: 10.1016/j.physletb.2014.06.047 Cite as:arXiv:1407.0011 [gr-qc] (or arXiv:1407.0011v1 [gr-qc] for this version) **Effects of dark energy on P-V criticality of charged AdS black holes Gu-Qiang Li STOP** Monitoring the orbits of stars around Sgr A* offers the possibility of detecting the precession of (e) their orbital planes due to frame dragging, of measuring the spin and quadrupole moment of the black hole, and of testing the no-hair theorem. Here authors investigate whether the deviations of stellar orbits from (e) test-particle trajectories due to (e) wind mass loss and tidal dissipation of (e) the orbital energy compromise such measurements. Authors find that the effects of stellar winds are, in general, negligible. On the other hand, for the most eccentric orbits (e>0.96) for (e) which an optical interferometer, such as GRAVITY, will detect (e&eb) orbital plane precession due to (e) frame dragging, the tidal dissipation of orbital energy occurs at (eb) timescales comparable to the timescale of precession due to (e) the quadrupole moment of the black hole. As a result, this non-conservative effect is a potential source of systematic uncertainty in testing the no-hair theorem with stellar orbits.Subjects:High Energy Astrophysical Phenomena (astro-ph.HE); Earth and Planetary Astrophysics (astro-ph.EP); General Relativity and Quantum Cosmology (gr-qc) DOI: 10.1088/0004-637X/777/1/57 Cite as: arXiv:1212.3342 [astro-ph.HE] (or arXiv:1212.3342v1 [astro-ph.HE] for this

version) **Deviation of Stellar Orbits from Test Particle Trajectories Around Sgr A* Due to Tides and Winds Dimitrios Psaltis (Arizona), Gongjie Li (Harvard), Abraham Loeb (Harvard) STOP** Authors perform a phase space analysis of a generalized modified gravity theory with non-minimal coupling between geometry and matter. Authors apply the dynamical system approach to (e&eb) this generalized model and find that in the cosmological context, different choices of (e) Lagrangian density will apparently result in (eb) different phases of the universe. By carefully choosing the variables, authors prove that there is an attractor solution to describe (eb)the late time accelerating universe when (e) the modified gravity is(=) chosen in a simple power-law form of (e) the curvature scalar. Authors further examine the temperature evolution based on (e) the thermodynamic understanding of the model. Confronting the model with (e&eb) temperature-redshift and (e&eb) supernova type Ia data sets; authors find that the non-minimally coupled theory of (e) gravity is (=) a viable model to describe (eb) the late time universe acceleration. Subjects: General Relativity and Quantum Cosmology (gr-qc); Cosmology and Nongalactic Astrophysics (astro-ph.CO); High Energy Physics - Theory (hep-th) Cite as: arXiv: 1512.09281 [gr-qc] (or arXiv:1512.09281v1 [gr-qc] for this version) **Dynamical analysis of modified gravity with non-minimal gravitational coupling to matter Rui An, Xiaodong Xu, Bin Wang, Yungui Gong STOP** The tensor-vector-scalar (TeVeS) model is considered a viable theory of gravity. It produces (eb) the Milgrom's modified Newtonian dynamics in (eb) the nonrelativistic weak field limit and is (=) free from ghosts. This model has been tested against (e&eb) various cosmological observations. Here authors investigate whether new observations such as the galaxy velocity power spectrum measured by (e) 6dF and the kinetic Sunyaev Zel'dovich effect power spectrum measured by (e) ACT/SPT can put further constraints on (eb) the TeVeS model. Furthermore, authors perform the test of TeVeS cosmology with (e&eb) a sterile neutrino by (e) confronting to (e) Planck data, and find that (eb) it is ruled out by (e) cosmic microwave background measurements from (e) the Planck mission. Subjects: Cosmology and Nongalactic Astrophysics (astro-ph.CO); General Relativity and Quantum Cosmology (gr-qc); High Energy

Physics - Theory (hep-th) Journal reference: Phys. Rev. D 92, 083505 (2015) DOI:10.1103/PhysRevD.92.083505 Cite as: arXiv:1412.4073 [astro-ph.CO] (or arXiv:1412.4073v3 [astro-ph.CO] for this version) **Testing the tensor-vector-scalar Theory with the latest cosmological observations Xiao-dong Xu, Bin Wang, Pengjie Zhang STOP** Using the anti–de Sitter/conformal field theory correspondence, authors relate the shear viscosity η of the finite-temperature N=4 supersymmetric Yang-Mills theory in the large N, strong-coupling regime with (e&eb) the absorption cross section of (e) low-energy gravitons by (e) a near-extremal black three-brane. **G. Policastro, D. T. Son, and A. O. Starinets** show that in the limit of zero frequency this cross section coincides with (=) the area of the horizon. From this result they find η=π8N2T3. They conjecture that for finite't Hooft coupling g2YMN the shear viscosity is (=) η=f(g2YMN)N2T3, where f(x) is a monotonic function that decreases from O(x−2ln−1(1/x)) at small x to π/8 when x→∞. Received 14 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevLett.87.081601 **Shear Viscosity of Strongly Coupled N=4 Supersymmetric Yang-Mills Plasma G. Policastro, D. T. Son, and A. O. Starinets Phys. Rev. Lett. 87, 081601 – Published 2 August 2001 STOP** Authors use the AdS/CFT correspondence to determine the rate of energy loss of a heavy quark moving through (e&eb) Script N = 4 SU (Nc) supersymmetric Yang-Mills plasma at large't Hooft coupling. Using the dual description of the quark as a classical string ending on a D7-brane, **Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1** use a complementary combination of analytic and numerical techniques to determine (eb) the friction coefficient as a function of quark mass. Provided strongly coupled Script N = 4 Yang-Mills plasma is (=) a good model for hot, strongly coupled QCD, results may be relevant for (e) charm and bottom physics at RHIC. **Energy loss of a heavy quark moving through Script N = 4 supersymmetric Yang-Mills plasma Christopher P. Herzog1, Andreas Karch1, Pavel Kovtun2, Can Kozcaz1 and Laurence G. Yaffe1 Published 11 July 2006 Journal of High Energy Physics, Volume 2006, JHEP07(2006) STOP Benjamin Svetitsky** calculates the classical drag and diffusion coefficients for (e) a charmed

quark propagating in the quark-gluon plasma. Both coefficients turn out rather large, so that (1) a charmed quark created when (e) the plasma is hot will be (=) stopped before propagating 1 fm and (2) subsequent diffusion will be fast. The first effect should serve to (e) increase the yield of J/ψ mesons in (eb) relativistic heavy-ion collisions, while the second should work in the opposite direction. In any case, the two effects should dominate (e) the dynamics of a $c\bar{c}$ pair. Received 16 November 1987 DOI: http://dx.doi.org/10.1103/PhysRevD.37.2484 **Diffusion of charmed quarks in the quark-gluon plasma Benjamin Svetitsky Phys. Rev. D 37, 2484 – Published 1 May 1988 STOP** Using holographic duality, authors present results for (e) both head-on and off-center collisions of Gaussian shock waves in (eb) strongly coupled N=4 supersymmetric Yang-Mills theory. The shock waves superficially resemble (e&eb) Lorentz contracted colliding protons. The collisions results in (eb) the formation of a plasma whose evolution is well described by (e) viscous hydrodynamics. The size of the produced droplet is R∼1/Teff where (e) Teff is the effective temperature, which is (=) the characteristic microscopic scale in strongly coupled plasma. These results demonstrate (eb) the applicability of hydrodynamics to microscopically small systems and bolster (eb) the notion that hydrodynamics can be applied to (e&eb) heavy-light ion collisions as well as some proton-proton collisions. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1601.01583 [hep-th] (or arXiv: 1601.01583v2 [hep-th] for this version) **How big are the smallest drops of quark-gluon plasma?**

All this consists of the act of consciousness alone", says Vasubandhu, "because unreal entities appear, just as a man with defective vision sees unreal hair or a moon, etc."He points to dreams as examples of purely subjective constructs which appear to be objective realities. The apparent reality dreams possess is not derived from any concrete, objective world, but merely from the idea of objectivity. While the Yogacara does not say that an idea has, for example, spatial attributes, it does have a form manifesting them. While he agrees with the Saiva idealist that

appearances have no independent existence apart from their appearing to consciousness, he maintains that for this reason they are unreal. The creativity of consciousness consists in its diversification in many modes having apparent externality; it is not a creation of objects. While the Kashmiri Saivite agrees that the world is pure conscious- ness alone, he maintains that it is such because it is a real creation of consciousness. The effect is essentially identical with the cause and shares in its reality. Matter and the entire universe are absolutely real, as 'congealed' (sty ana) or 'contracted' (samkucita) forms of consciousness. "This God of consciousness", writes Ksemaraja, "generates the universe and its form is a condensation of His own essence (rasa)" m By boiling sugarcane juice it condenses to form treacle, brown sugar and candy which retains its sweetness. Similarly, consciousness abides unchanged even though it assumes the concrete material form of the five gross elements. The same reality thus abides equally in gross and subtle forms. Consequently no object is totally insentient. Even stones bear a trace (vasana) of consciousness, although it is not clearly apparent because it is not associated with the vital breath (prana) and other components of a psycho-physical organism. Somananda goes so far as to affirm that physical objects, far from being insentient, can only exist insofar as they are aware of themselves as existing. The jar performs its function because it knows itself to be its agent. Indeed, all things are pervaded by consciousness and at one with it and hence share in its omniscience. Thus, Siva, Who perceives Himself in the form of physical objects, is the one ultimate reality. "The jar knows because it is of my nature", writes Somananda, "and I know it because I am of the jar's nature. I know because I am of Sadasiva's nature and He knows because He is of my nature; Yajnadatta [knows] because he is of Siva's nature and Siva [knows] because He is of Yajnadatta's nature". Integral Monism of Kashmiri Saivism Everything in this sense is directly perceived by absolute conscious- ness, and this direct perception (pratyaksa) unifies the knowable into a single, undivided whole. This is the central concept behind a doctrine originally expounded by Narasimha called 'the non-dualism of direct perception' (pratyaksadvaita). This states that consciousness is essentially perceptive and that its perception of all things operates throughout the universe. Insofar

as phenomena are clearly evident (sphufa) to us, everything is directly perceived by absolute consciousness, with which our individual consciousness is identical. This direct perception unfolds everywhere; the one true reality, it is alone and without companion or rival (nihsapatna). Even though it remains one, it can, by its very nature, perceive distinctions (bheda) between one entity and another, without this engendering any division within it. We distinguish between two entities in empirical terms on the basis of their mutual exclusion (anyonyabhava). The relative distinction {bheda) between them is essentially the perceived difference between their respective characteristics. Despite this difference they are united within the purview of a single cognition insofar as they are equally both manifest appearances. This cognition is the undivided essence (rasa) or 'own nature' (svabhava) of both. Encompassed by the 'fire of consciousness', there is no essential difference between them. Just as when an emerald and ruby reflect each other's light, the ruby is reddish-green and the emerald greenish-red, similarly everything is connected with everything else as part of the single variegated (vicitra) cognition of absolute consciousness. Mahesvarananda writes: The Supreme Lord's unique state of emotivity (asadharanabhava) is the outpouring of pure Being (mahasatta). It is manifest as the brilliance (sphuratta) of the universe which, if we ponder deeply, [is realized to be] the single flavour (ekarasa) of the essence of Beauty which is the vibration of the bliss of one's own nature. In this way all things are in reality one although divided from the one another sharing as they do the 'single flavour' (ekarasa) of the pure vibration of consciousness. Kashmiri Saiva Realism Kashmiri Saivism as a whole has been variously called a form of 'realistic idealism', 'monistic idealism', ‘idealistic monism' and 'concrete monism'. It is easy to understand why Kashmiri Saivism is The Doctrine of Vibration said to be 'idealistic' and 'monistic', but in what sense is it also 'realistic*? The answer to this question is of no small importance in trying to under- stand the central idea behind its metaphysics and the fundamental importance of the concept of Spanda, in this seemingly impossible marriage between monistic idealism and pluralistic realism. The Kashmiri Saiva approach understands the world to be a symbol of the absolute, that is, as the manner in which it presents itself to us. Again we can

contrast this view with that of the Advaita Vedanta. The Advaita Vedanta understands the world to be an expression of the absolute insofar as it exists by virtue of the absolute's Being. Being is understood to be the real unity which underlies empirically manifest separateness and as such is never empirically manifest. It is only transcendentally actual as 'being-in-itself. The Kashmiri Saiva position represents, in a sense, a reversal of this point of view. The nature of the absolute, and also that of Being, is onceived as an eternal becoming (satatodita), a dynamic flux or Spanda, 'the agency of the act of being'. It is dentified with the concrete actuality of the fact of appearing, not passive unmanifest Being. Appearance (abhasa) alone is real.ppearing (prakasamanatva) is equivalent to the fact of being (astitva). Ksemaraja writes in his commentary on the Stanzas on Vibration: Indeed, all things are manifest because they are nothing but manifestation. The point being that nothing is manifest apart from manifestation The absolutely unmanifest, from this point of view, can have as little existence as the space in a lattice window of a sky-palace. Nay, even less, because even that space can appear as an imagined mage manifest within consciousness. Everything is real according to the manner in which it appears. Even an illusion is in this sense real, insofar as it appears and is known in the manner in which it appears. The empirical and the real are identical categories of thought. As Abhinava says: Thus this is the supreme doctrine (upanisatf), namely that, when- ever and in whatever form [an entity] appears, that then is its particular nature. Perhaps at this stage a brief comparison with Heidegger's ideas might prove to be enlightening and not altogether out of place. According to Heidegger's phenomenology of Being, reality is intelligible in a two-fold Integral Monism of Kashmiri Saivism manner as 'phenomenon' and 'logos'. Heidegger defines what he means by 'phenomenon' as: "that-which-shows-itself. The manifest . . . phenomena are then the collection of that which lies open in broad daylight or can be brought to the light of day — what the Greeks at times implicitly identified as 'ta onta' (the things-which-are)".

**ABSOLUTE STABILITY WITH A DISTRIBUTED DELAY DUE**

To debit consummation

It is shown elsewhere (Nunney, 1985a) that the characteristic equation of the model (3) is given by

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 e^{-\tau_9\lambda_9} + X_9 + Y_9 = 0 \qquad (A1.1)$$

Where $V_9 = dB_9/dR_9$ $U_9 = dD_9/dR_9 + N_9 dG_9/dR_9$ $X_9 = N_9 G_9 dF_9/dR_9$ and $Y_9 = -N_9 G_9 dM_9/dR_9$ (evaluated at the equilibrium). Note that $U_9 X_9$ and $Y_9$ are each nonnegative.

We now assume that the delay is not of fixed length T, but varies between debit –credit. We can define the probability that the delay is of length $T_9$ as $(T_9)$ . (It follows that $f(T_9)$ 0 for all $T$, and $\int_0^\infty f(T_9)dt = 9.$) Incorporating this variation in Eq.(A1.1) gives

$$\lambda_9^2 + U_9\lambda_9 - V_9\lambda_9 \int_0^\infty e^{-\tau_9\lambda_9} f(T_9) + X_9 + Y_9 = 0. \qquad (A1.2)$$

In order to investigate the stability properties of (A1.2) we employ the method of D-partitions (El'sgolt's and Norkin, 1973) in which parameter space is divided by surfaces which define at least one root with a zero real

**Paul M. Chesler, Krishna Rajagopal** define the rate of energy loss dEjet/dx and (e&eb) the jet opening angle in a straightforward fashion directly in the gauge theory before (e) calculating both holographically, in (eb) the dual gravitational description. In this way, authors rederive the previously known result for (e) dEjet/dx without (e) the need to introduce (e&eb) a finite slab of plasma. They obtain a striking relationship between the initial opening angle of the jet, which is to say the opening angle that it would have had if it had found itself in vacuum instead of in plasma, and (e&eb) the thermalization distance of the jet. Via this relationship, they show that N=4 SYM jets with (e&eb) any initial energy that have the same initial opening angle and the same trajectory through (e&eb) the plasma experience the same fractional energy loss. They also provide an expansion that describes how the opening angle of the N=4 SYM jets increases slowly as they lose (e) energy, over (e&eb) the fraction of their lifetime when (e) their fractional energy loss is (=) not yet large. Authors close by looking ahead toward potential qualitative lessons from (e) results for QCD jets produced in (eb) heavy collisions and propagating through (e&eb) quark-gluon plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1511.07567 [hep-th] (or arXiv:1511.07567v2 [hep-th] for this version) **On the Evolution of Jet Energy and Opening Angle in Strongly Coupled Plasma Paul M. Chesler, Krishna**

**Rajagopal STOP Paul M. Chesler, Niki Kilbertus, Wilke van der Schee study** the collision of planar shock waves in (eb) AdS5 as a function of (e) shock profile. In the dual field theory the shock waves describe (eb) planar sheets of energy whose collision results in (eb) the formation of a plasma which behaves (e&eb) hydrodynamically at late times. They find that the post-collision stress tensor near (eb) the light cone exhibits (eb) transient non-universal behavior which depends on (=) both the shock width and (e&eb) the precise functional form of (e) the shock profile. However, over a large range of shock widths, including those which yield (eb) qualitative different behavior near (eb) the future light cone, and for (e) different shock profiles, authors find universal behavior in (eb) the subsequent hydrodynamic evolution. Additionally, authors compute the rapidity distribution of (e) produced particles and find it to be well described by (e) a Gaussian. Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) DOI: 10.1007/JHEP11 (2015)135 Cite as: arXiv: 1507.02548 [hep-th] (or arXiv: 1507.02548v2 [hep-th] for this version) **Universal hydrodynamic flow in holographic planar shock collisions Paul M. Chesler, Niki Kilbertus, Wilke van der Schee STOP** Using numerical holography, authors study the collision, at non-zero impact parameter, of bounded, localized distributions of energy density chosen to (e) mimic relativistic heavy ion collisions, in (eb) strongly coupled N = 4 supersymmetric Yang-Mills theory. Both longitudinal and transverse dynamics in the dual field theory are (=) properly described. Using the gravitational description we solve 5D Einstein equations, without dimensionality reducing symmetry restrictions, to find the asymptotically anti-de Sitter spacetime geometry. Implications of **Paul M. Chesler, Laurence G. Yaffe** results on the understanding of early stages of heavy ion collisions, including the development of transverse radial flow, are discussed. Comments: Enlarged JHEP version, 14 pages Subjects: High Energy Physics - Theory (hep-th); General Relativity and Quantum Cosmology (gr-qc); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Cite as: arXiv: 1501.04644 [hep-th] (or arXiv:1501.04644v2 [hep-th] for this version) **Holography and off-center collisions of localized shock waves Paul M. Chesler, Laurence G. Yaffe STOP Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** study the dynamic after a smooth quench across a continuous transition from the disordered phase to the ordered phase. Based on scaling ideas, linear response and (e&eb) the spectrum of unstable modes, authors develop a theoretical framework, valid for (e) any second order phase transition, for the early-time evolution of the condensate in the broken phase. Their analysis unveils (eb) a novel

period of non-adiabatic evolution after the system passes through (e&eb) the phase transition, where (e) a parametrically large amount of coarsening occurs before (e) a well-defined condensate forms. Formalism predicts (eb) a rate of defect formation parametrically smaller than (eb) the Kibble-Zurek prediction and yields (eb) a criterion for (e) the break-down of Kibble-Zurek scaling for sufficiently fast quenches. **Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu** numerically test formalism for (e) a thermal quench in (eb) a 2 + 1 dimensional holographic superfluid. These findings, of (e) direct relevance in a broad range of fields including cold atom, condensed matter, statistical mechanism and cosmology, ar (=)e an important step towards a more quantitative understanding of dynamical phase transitions. Subjects: High Energy Physics - Theory (hep-th); Quantum Gases (cond-mat.quant-gas); Statistical Mechanics (cond-mat.stat-mech) Journal reference: Phys. Rev. X 5, 021015 (2015) DOI: 10.1103/PhysRevX.5.021015 Cite as: arXiv:1407.1862 [hep-th] (or arXiv:1407.1862v2 [hep-th] for this version) **Defect formation beyond Kibble-Zurek mechanism and holography Paul M. Chesler, Antonio M. Garcia-Garcia, Hong Liu STOP Paul M. Chesler, Krishna Rajagopal** present calculations in which an energetic light quark shoots through (e&eb) a finite slab of strongly coupled N=4 supersymmetric Yang-Mills (SYM) plasma, with thickness L, focussing on what comes out on (eb) the other side. Authors find that even when the "jets" that emerge from (e) the plasma have lost a substantial fraction of (e) their energy they look in almost all respects like "jets" in (eb) vacuum with (e&eb) the same reduced energy. The one possible exception is (=) that the opening angle of the "jet" is (=) larger after passage through (e&eb) the slab of plasma than before. Along the way, authors obtain (eb) a fully geometric characterization of energy loss in (eb) the strongly coupled plasma and show that (eb) dE out/dL∝L2/x2stop−L2−−−−−−−−√, where Eout is the energy of the "jet" that emerges from (e) the slab of plasma and x stop is (=) the (previously known) stopping distance for (e) the light quark in an infinite volume of (e) plasma. Subjects: High Energy Physics - Theory (hep-th); High Energy Physics - Phenomenology (hep-ph); Nuclear Theory (nucl-th) Journal reference: Phys. Rev. D 90, 025033 (2014) DOI: 10.1103/PhysRevD.90.025033 Cite as: arXiv: 1402.6756 [hep-th] (or arXiv: 1402.6756v1 [hep-th] for this version) **Jet quenching in strongly coupled plasma Paul M. Chesler, Krishna Rajagopal STOP** Quasinormal modes are (=) eigenmodes of dissipative systems. Perturbations of classical gravitational backgrounds involving (e&eb) black holes or branes naturally lead to (eb) quasinormal modes. The analysis and classification of (e) the quasinormal spectra require solving (e) non-Hermitian eigenvalue problems for (e) the associated linear

differential equations. Within the recently developed gauge-gravity duality, these modes serve as (=) an important tool for (e) determining the near-equilibrium properties of (e) strongly coupled quantum field theories, in particular their transport coefficients, such as (=) viscosity, conductivity and diffusion constants. In astrophysics, the detection (e&eb) of quasinormal modes in gravitational wave experiments would allow (eb) precise measurements of (e) the mass and spin of black holes as well as new tests of general relativity. This review is meant as an introduction to the subject, with a focus on the recent developments in the field. **Quasinormal modes of black holes and black branes Emanuele Berti1,2, Vitor Cardoso1,3 and Andrei O Starinets4 Published 24 July 2009 • 2009 IOP Publishing Ltd Classical and Quantum Gravity, Volume 26, Number 16 STOP Gary T. Horowitz and Veronika E. Hubeny** investigate the decay of a scalar field outside (e) a Schwarzschild anti–de Sitter black hole. This is determined by computing (e&eb) the complex frequencies associated with (e&eb) quasinormal modes. There are qualitative differences from (e) the asymptotically flat case, even in (eb) the limit of small black holes. In particular, for a given angular dependence, the decay is (=) always exponential—there are (=) no power law tails at late times. In terms of the AdS-CFT correspondence, a large black hole corresponds to (e&eb) an approximately thermal state in the field theory, and the decay of the scalar field corresponds to (e&eb) the decay of a perturbation of this state. Thus one obtains the time scale for the approach to (e) thermal equilibrium. **Gary T. Horowitz and Veronika E. Hubeny** compute (e&eb) these time scales for the strongly coupled field theories in three, four, and six dimensions, which are (=)dual to string theory in asymptotically AdS spacetimes. Received 27 September 1999 DOI: http://dx.doi.org/10.1103/PhysRevD.62.024027 **Quasinormal modes of AdS black holes and the approach to thermal equilibrium Gary T. Horowitz and Veronika E. Hubeny Phys. Rev. D 62, 024027 – Published 27 June 2000 STOP** Gravitational waves emitted by (e) perturbed black holes or relativistic stars are dominated by (e) `quasinormal ringing', damped oscillations at (eb) single frequencies which are (=) characteristic of the underlying system. These quasinormal modes have been studied for (e) a long time, often with the intent of describing the time evolution of (e) a perturbation in terms of (e&eb) these modes in a way very similar to a normal-mode analysis. In this review, authors summarize how quasinormal modes are defined and computed. **Hans-Peter Nollert** finds why they have been regarded as closely analogous to normal modes, and discovers why they are actually quite different. Authors also discuss how quasinormal modes can be used in (eb) the analysis of a gravitational wave signal, such as will hopefully be detected in (e&eb) the near future. **Quasinormal modes: the characteristic `sound' of black**

**holes and neutron stars Hans-Peter Nollert STOP** Perturbations of black holes, initially considered in the context of possible observations of astrophysical effects, have been studied for the past 10 years in string theory, brane-world models, and quantum gravity. Through the famous gauge/gravity duality, proper oscillations of (e) perturbed black holes, called quasinormal modes, allow (eb) for the description of (e) the hydrodynamic regime in (eb) the dual finite temperature field theory at (eb) strong coupling, which can be used to (e) predict the behavior of quark-gluon plasmas in (eb) the nonperturbative regime. On the other hand, the brane-world scenarios assume (eb) the existence of extra dimensions in nature, so that (e) multidimensional black holes can be formed in (eb) a laboratory experiment. All this stimulated active research in the field of perturbations of (e) higher-dimensional black holes and branes during recent years. In this review recent achievements on various aspects of black hole perturbations are discussed such as decoupling of variables in the perturbation equations, quasinormal modes (with special emphasis on various numerical and analytical methods of calculations), late-time tails, (e&eb) gravitational stability, (e&eb) anti–de Sitter/conformal field theory interpretation of (e) quasinormal modes, and holographic superconductors. Authors also touch on state-of-the-art observational possibilities for (e) detecting quasinormal modes of black holes. DOI: http://dx.doi.org/10.1103/RevModPhys.83.793 © 2011 American Physical Society **Quasinormal modes of black holes: From astrophysics to string theory R. A. Konoplya and Alexander Zhidenko Rev. Mod. Phys. 83, 793 – Published 11 July 2011 STOP Vitor Cardoso and José P. S. Lemos** study the quasinormal modes (QNM) of electromagnetic and gravitational perturbations of a Schwarzschild black hole in (eb) an asymptotically anti–de Sitter (AdS) spacetime. Some of the electromagnetic modes do not (e) oscillate; they only decay, since they have (e) pure imaginary frequencies. The gravitational modes show (eb) peculiar features: the odd and even gravitational perturbations no longer have (e) the same characteristic quasinormal frequencies. There is a special mode for odd perturbations whose behavior differs completely from (e) the usual one in scalar and electromagnetic perturbations in AdS spacetime, but has a similar behavior to the Schwarzschild black hole in an asymptotically flat spacetime: the imaginary part of the frequency goes as 1/r+, where r+ is the horizon radius. **Vitor Cardoso and José P. S. Lemos** also investigate the small black-hole limit showing that the imaginary part of the frequency goes as r2+. These results are important to the AdS/CFT conjecture since, according to it, the QNM's describe the approach to equilibrium in the conformal field theory. Received 12 April 2001 DOI: http://dx.doi.org/10.1103/PhysRevD.64.084017 **STOP Gianluca Calcagni** studies the ultraviolet complete non-

relativistic theory recently proposed by Hořava. After introducing a Lifshitz scalar for a general background, authors analyze the cosmology of (e) the model in Lorentzian and Euclidean signature. Vacuum solutions are found and it is argued the existence of (e) non-singular bouncing profiles. They find a general qualitative agreement with both the picture of Causal Dynamical Triangulations and (e&eb) Quantum Einstein Gravity. However, inflation driven by (e) a Lifshitz scalar field on a classical background might not (e) produce a scale-invariant spectrum when (e) the principle of detailed balance is assumed. **Cosmology of the Lifshitz universe Gianluca Calcagni Published 25 September 2009 • Journal of High Energy Physics, Volume 2009, JHEP09 (2009) STOP Petr Hořavaa, , Edward Wittenb** propose that the ten-dimensional E8 × E8 heterotic string is related to an eleven-dimensional theory on the orbifold Full-size image (<1 K) in the same way that the Type IIA string in ten dimensions is related to Full-size image (<1 K). This in particular determines the strong coupling behavior of the ten-dimensional E8 × E8 theory. It also leads to a plausible scenario whereby duality between SO (32) heterotic and Type I superstrings follows from the classical symmetries of the eleven-dimensional world, just as (=) the Full-size image (<1 K) duality of the ten-dimensional Type IIB theory follows from eleven-dimensional diffeomorphism invariance. **Heterotic and Type I string dynamics from eleven dimensions Petr Hořavaa, , Edward Wittenb, Nuclear Physics B Volume 460, Issue 3, 12 February 1996, Pages 506–524 STOP Petr Hořava1** proposes a quantum theory of membranes designed such that (e) the ground-state wavefunction of the membrane with (e&eb) compact spatial topology Σh reproduces (eb) the partition function of the bosonic string on (e&eb) worldsheet Σh. The construction involves (e&eb) worldvolume matter at (eb) quantum criticality, described in the simplest case by (e) Lifshitz scalars with (e&eb) dynamical critical exponent z = 2. This matter system must be coupled to (e&eb) a novel theory of worldvolume gravity, also exhibiting (eb) quantum criticality with z = 2. Authors first construct such a nonrelativistic ``gravity at a Lifshitz point" with z = 2 in D+1 spacetime dimensions, and then specialize to (e&eb) the critical case of D = 2 suitable for the membrane worldvolume. They also show that (eb) in the second-quantized framework, the string partition function is reproduced if (e) the spacetime ground state takes the form of (e) a Bose-Einstein condensate of membranes in (eb) their first-quantized ground states, correlated across all general. **Membranes at quantum criticality Petr Hořava1, 2 Published 4 March 2009 • Journal of High Energy Physics, Volume 2009, JHEP03 (2009)STOP Petr Hořava** presents a candidate quantum field theory of gravity with (e&eb) dynamical critical exponent equal to z=3 in the UV. (As in condensed-matter systems, z measures (eb)

the degree of anisotropy between space and time.) This theory, which at short distances describes (eb) interacting nonrelativistic gravitons, is (=) power-counting renormalizable in (eb) 3+1 dimension. When restricted to satisfy the condition of detailed balance,(eb) this theory is intimately related to (e&eb) topologically massive gravity in three dimensions, and the geometry of (e) the **Cotton tensor**. At long distances, this theory flows (e&eb) naturally to the relativistic value z=1, and could therefore serve as (=) a possible candidate for (e) a UV completion of Einstein's general relativity or an infrared modification thereof. The effective speed of light, the Newton constant and the cosmological constant all emerge from (e) relevant deformations of the deeply nonrelativistic z=3 theory at short distances. **Quantum gravity at a Lifshitz point Petr Hořava Phys. Rev. D 79, 084008 – Published 6 April 2009STOP** Black holes appear to conform to (e) a very straightforward generalisation of standard laboratory thermodynamics. This generalised theory is examined in detail, and some concrete results are presented. The thermodynamic connection is based on (e) Hawking's application of quantum theory to (e) black holes, and the quantum aspects are described in detail from (e) several standpoints, both heuristic and otherwise. The precise mechanism by which the black hole produces thermal radiation, its nature and origin, and the energetics of back-reaction on the hole are reviewed. The thermal states of (e) quantum holes are also treated using (e) the theory of thermal Green functions, and the entropy of the hole is shown to be related to (e&eb) the loss of information about the quantum states hidden (e) behind the event horizon. Some related topics such as accelerated mirrors and (e&eb) observers in Minkowski space; (e&eb) super-radiance from (e) rotating holes and the thermodynamics of general self-gravitating systems are also briefly discussed. **Thermodynamics of black holes P C W Davies Reports on Progress in Physics, Volume 41, and Number STOP D. Blasa, O. Pujolàsb and S. Sibiryakova,c** address the consistency of Hořava's proposal for a theory of quantum gravity from the low-energy perspective. They uncover the additional scalar degree of freedom arising from (e) the explicit breaking of the general covariance and study its properties. The analysis is performed both in (eb) the original formulation of the theory and in (eb) the Stückelberg picture. A peculiarity of the new mode is that it satisfies (eb) an equation of motion that is of (e) first order in time derivatives. At linear level the mode is manifest only around (e&eb) spatially inhomogeneous and time-dependent backgrounds. Authors find two serious problems associated with (e&eb) this mode. First, the mode develops (eb) very fast exponential instabilities at (eb) short distances. Second, it becomes (eb) strongly coupled at (eb) an extremely low cutoff scale. They also discuss the ``projectable'' version of (e) Hořava's proposal and argue that this

version can be understood as (=) a certain limit of the ghost condensate model. The theory is still problematic since (e) the additional field generically forms (eb) caustics and, again, has a very low strong coupling scale. Authors clarify some subtleties that arise in (eb) the application of the Stückelberg formalism to Hořava's model due to its non-relativistic nature. **On the extra mode and inconsistency of Hořava gravity D. Blasa, O. Pujolàsb and S. Sibiryakova,c Published 12 October 2009 • Journal of High Energy Physics, Volume 2009, JHEP10(2009) STOP Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** identify signatures of chaos in the dynamics of discord in a multiqubit system collectively modeled as (=) a quantum kicked top. Evolution of discord between any two qubits is quasiperiodic in regular regions, while (e) in chaotic regions the quasiperiodicity is lost. As the initial wave function is varied from (e) the regular regions to the chaotic sea, a contour plot of the time-averaged discord remarkably reproduces (eb) the structures of the classical stroboscopic map. **Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose** also find surprisingly **opposite behavior** of two-qubit discord versus (e&eb) entanglement of the two qubits as measured by the concurrence. Results provide evidence of signatures of (e) chaos in dynamically generated discord. http://dx.doi.org/10.1103/PhysRevE.91.032906©2015 American Physical Society **Signatures of chaos in the dynamics of quantum discord Vaibhav Madhok, Vibhu Gupta, Denis-Alexandre Trottier, and Shohini Ghose Phys. Rev. E 91, 032906 – Published 10 March 2015 STOP.** How do closed quantum many-body systems driven out of (e) equilibrium eventually achieve equilibration? And how do these systems thermalize, given that they comprise so many degrees of freedom? Progress in answering these—and related—questions has accelerated in recent years—a trend that can be partially attributed to success with experiments performing quantum simulations using ultracold atoms and trapped ions. **J. Eisert, M. Friesdorf & C. Gogolin** provide (eb) an overview of this progress, specifically in studies probing dynamical equilibration and (e&eb) thermalization of systems driven out of (e) equilibrium by quenches, ramps and periodic driving. In doing so, they also address (e&eb) topics such as the eigenstate thermalization hypothesis, typicality, transport, many-body localization and universality near phase transitions, as well as future prospects for quantum simulation. **Quantum many-body systems out of equilibrium J. Eisert, M. Friesdorf& C. Gogolin Nature Physics 11, 124–130 (2015) doi: 10.1038/nphys3215 Received 11 August 2014 Accepted 03 December 2014 published online 03 February 2015 STOP.** Quantum teleportation is (=) one of the most important protocols in quantum information. By exploiting the physical resource of

entanglement, quantum teleportation serves as a key primitive in a variety of quantum information tasks and represents an important building block for quantum technologies, with a pivotal role in the continuing progress of quantum communication, quantum computing and quantum networks. **Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein** review the basic theoretical ideas behind quantum teleportation and its variant protocols. They focus on the main experiments, together with (e&eb) the technical advantages and disadvantages associated with (e&eb) the use of the various technologies, from photonic qubits and (e&eb) optical modes to atomic ensembles, trapped atoms, and (e&eb) solid-state systems. Analysing the current state-of-the-art, authors finish by discussing open issues, challenges and potential future implementations. Cite as: arXiv: 1505.07831 [quant-ph] (or arXiv: 1505.07831v1 [quant-ph] for this version) **Advances in Quantum Teleportation Stefano Pirandola, Jens Eisert, Christian Weedbrook, Akira Furusawa, and Samuel L. Braunstein (Submitted on 28 May 2015) STOP. Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are channels for which the coherent information is zero, but which nonetheless have capacity. Transmitting data reliably over (e&eb) noisy communication channels is one of the most important applications of (e) information theory, and is well understood for (e) channels modeled by classical physics. However, when quantum effects are involved, (eb) one do not know how to compute channel capacities. This is because the formula for (e) the quantum capacity involves (e&eb) maximizing the coherent information over (e&eb) an unbounded number of (e) channel uses. In fact, entanglement across channel uses can even increase the coherent information from zero to non-zero. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** study the number of channel uses necessary to detect positive coherent information. In all previous known examples, two channel uses already sufficed. It might be that only a finite number of channel uses are always sufficient. **Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk** show that this is not the case: for any number of uses, there are (=) channels for which the coherent information is (=) zero, but which nonetheless have (e) capacity. **Unbounded number of channel uses may be required to detect quantum capacity Toby Cubitt, David Elkouss, William Matthews, Maris Ozols, David Pérez-García & Sergii Strelchuk Nature Communications 6, Article number: 6739 doi:10.1038/ncomms7739 Published 31 March 2015 STOP.** Characterizing the behaviour of strongly coupled

quantum systems out of (e) equilibrium is a cardinal challenge for (e) both theory and experiment. With diverse applications ranging from the dynamics of the quark–gluon plasma to transport in novel states of quantum matter, establishing universal results and organizing principles out of equilibrium is crucial. **M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm** present a universal description of energy transport between quantum critical heat baths in arbitrary dimension. Current-carrying non-equilibrium steady state (NESS) is a Lorentz-boosted thermal state. In the context of gauge/gravity duality this reveals an intimate correspondence between far-from-equilibrium transport and black hole uniqueness theorems. Authors provide analytical expressions for (e) the energy current and the generating function of (e) energy current fluctuations, together with (e&eb) predictions for experiment. **Energy flow in (eb) quantum critical systems far from (e) equilibrium M. J. Bhaseen, Benjamin Doyon, Andrew Lucas& Koenraad Schalm Nature Physics 11, 509–514 (2015) doi: 10.1038/nphys3320: Published online 04 May 2015 STOP.** Properties of systems near quantum critical points (QCPs) have been studied extensively. A QCP is (=) a point across which the symmetry of (e) the ground state of a quantum system changes in (eb) a fundamental way; such a point can be accessed by (e) changing some parameter, say λ, in the Hamiltonian governing the system. The change in the ground state across (e&eb) a QCP is mediated by (e) quantum fluctuations. Unlike conventional thermal critical points, thermal fluctuations do not (e) play a crucial role in such transitions. Similar to its thermal counterparts, the low-energy physics near (e&eb) a QCP is associated with a number of critical exponents which characterize the universality class of such a transition. Among these exponents, the dynamical critical exponent z provides (eb) the signature of the relative scaling of space and time at (eb) the transition and has (e) no counterpart in thermal phase transitions. The other exponent which is going to be important for (e) the purpose of this review is (=) the well-known correlation length exponent ν. These exponents are formally defined as follows. As we approach the critical point at $\lambda=\lambda c$, the correlation length diverges as $\xi\sim|\lambda-\lambda c|-\nu$, while (e) the gap between the ground state and first excited state vanishes as (e) $\Delta E\sim\xi-z\sim|\lambda-\lambda c|z\nu$. Exactly at the critical point $\lambda=\lambda c$, the energy of the low-lying excitations vanishes at (eb) some wave number k 0 as $\omega\sim|k-k0|z$. The critical exponents are independent of (e) the details of the microscopic Hamiltonian; they depend only on a few parameters such as the dimensionality of the system and the symmetry of the order parameter. These features render (eb) the low-energy equilibrium physics of a quantum system near a QCP truly universal. **Quantum Quenching, Annealing and Computation Lecture Notes in Physics Volume 802, 2010, pp 21-56 Date: 08 Mar 2010 Non-**

**equilibrium Dynamics of Quantum Systems: Order Parameter Evolution, Defect Generation, and Qubit Transfer S. Mondal, D. Sen, K. Sengupta STOP Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl** propose a class of qubit networks that admit (e) the perfect state transfer of (e) any quantum state in a fixed period of time. Unlike many other schemes for (e) quantum computation and communication, these networks do not (e) require qubit couplings to be switched on and off. When restricted to (e) N-qubit spin networks of identical qubit couplings, they show that (eb) 2log3N is the maximal perfect communication distance for (e) hypercube geometries. Moreover, if one allows fixed but different couplings between the qubits, then (eb) perfect state transfer can be achieved over (e&eb) arbitrarily long distances in a linear chain. **Perfect State Transfer in Quantum Spin Networks Matthias Christandl, Nilanjana Datta, Artur Ekert, and Andrew J. Landahl Phys. Rev. Lett. 92, 187902 – Published 4 May 2004 STOP.** As with classical information processing, a quantum information processor requires (e) bits (qubits) that can be independently addressed (e&eb) and read out, long-term memory elements to store (e) arbitrary quantum states1, 2, and the ability to transfer (e&eb) quantum information through (e&eb) a coherent communication bus accessible to (e) a large number of qubits3, 4. Superconducting qubits made with (e&eb) scalable microfabrication techniques are (=) a promising candidate for the realization of (e) a large-scale quantum information processor5, 6, 7, 8, and 9. (For references please see the paper). Although these systems have successfully passed tests of (e) coherent coupling for (e) up to four qubits10, 11, 12, 13, communication of individual quantum states between (e&eb) superconducting qubits via (e&eb) a quantum bus has not yet been realized. Here, authors perform an experiment demonstrating the ability to coherently transfer quantum states between (e&eb) two superconducting Josephson phase qubits through a quantum bus. This quantum bus is (=) a resonant cavity formed by an open-ended superconducting transmission line of length 7 mm. After preparing an initial quantum state with (e&eb) the first qubit, this quantum information is transferred and stored (e&eb) as a nonclassical photon state of (e) the resonant cavity, then retrieved later by (e) the second qubit connected to (e&eb) the opposite end of the cavity. Beyond simple state transfer, these results suggest (eb) that a high-quality-factor superconducting cavity could also function as (=) a useful short-term memory element.

**Model has been applied for the following results:**

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| **Searches for Sterile Neutrinos with the IceCube Detector M. G. Aartsen et al. (IceCube Collaboration) Phys.**<br><br>**Rev. Lett 117, 071801 – Published 8 August 2016 Physics logo See Viewpoint: Hunting the Sterile Neutrino** |
| **(1) The IceCube neutrino telescope at the South Pole has measured the atmospheric muon neutrino spectrum as (=) a function of zenith angle and (e&eb) energy in the approximate 320 GeV to 20 TeV range, to search for (e) the oscillation signatures of (e) light sterile neutrinos.** |
| **(2) No evidence for anomalous νμ or ν¯μ disappearance is observed in (eb) either of two independently developed analyses, each using one year of (e) atmospheric neutrino data.** |
| **(3) New exclusion limits are placed on (e&eb) the parameter space of (e) the 3+1 model, in which muon antineutrinos experience (e) a strong Mikheyev-Smirnov-Wolfenstein-resonant oscillation.** |
| **(4) The exclusion limits extend to (e) sin22024≤0.02 at Δm2∼0.3 eV2 at the 90% confidence level.** |
| **(5) The allowed region from (e) global analysis of appearance experiments, including (e) LSND and MiniBooNE, is excluded at approximately the 99% confidence level for (e) the global best-fit value of \|Ue4\|2. Received 5 May 2016 DOI:http://dx.doi.org/10.1103/PhysRevLett.117.071801 © 2016 American Physical Society** |
| **Synopsis: Searching for Majorana Neutrinos August 16, 2016The KamLAND-Zen collaboration has run the most sensitive search to date for a radioactive decay that could reveal whether neutrinos are Majorana fermions. Synopsis figureKamLAND-Zen Collaboration** |
| **(6) Are neutrinos and (=) antineutrinos the same particles, as hypothesized by (e) physicist Ettore Majorana?** |
| **(7) If true, this "Majorana" nature of neutrinos could explain (eb) why neutrinos have (e) mass and why the Universe has (e) more matter than antimatter.** |
| **(8) Using an underground neutrino detector in Japan, the KamLAND-Zen collaboration has carried out the most sensitive search to date for a process called neutrinoless double-beta decay, which could confirm (eb) Majorana's hypothesis.** |
| **(9) The results set stringent constraints on the possible properties of (e) Majorana neutrinos.** |
| **(10) Double-beta decay is (=) a nuclear process in which two neutrons turn into (e&eb) two protons (or vice versa), emitting (eb) two antineutrinos.** |
| **(11) If neutrinos are Majorana fermions, (eb) the two antineutrinos could annihilate each other, and some double-beta decays would emit (eb) no antineutrinos.** |
| **(12) However, neutrinoless decay is believed to be (=) extremely rare. Searches of the process using (e) large amounts of the 35 naturally occurring (eb) isotopes that can undergo (e&eb) double-beta decay (such as xenon-136) have so far come up empty.** |

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| **(13) The KamLAND-Zen collaboration boosted their experimental sensitivity to (e) neutrinoless double-beta decays through (e&eb) a combination of factors.** |
| **(14) The researchers deployed clean detectors with (e&eb) very low background noise.** |
| **(15) And they used an unprecedented amount of xenon-136, which they purified to (e) remove any radioactive contaminants that could produce (eb) unwanted signals in the detectors.** |
| **(16) Combining data from several years, they improved (eb) the limit on the probability of neutrinoless double-beta decay by (e) sixfold compared to previous searches.** |
| **(17) This probability corresponds to (e&eb) a decay lifetime whose staggering value exceeds (e) 1026 years.** |
| **(18) This measurement allowed (eb) the researchers to set (eb) the most stringent upper limit on (e&eb) the possible mass of Majorana neutrinos (the particle must be lighter than 61–165 meV). This research is published in Physical Review Letters. –Matteo Rini Matteo Rini is the Deputy Editor of Physics.** |
| **How to Reach a Thousand-Second in-Plane Polarization Lifetime with 0.97−GeV/c Deuterons in a Storage Ring G. Guidoboni et al. (JEDI Collaboration) Phys. Rev. Lett. 117, 054801 – Published 28 July 2016 Physics logo See Synopsis: Deuterons Spin together for a Thousand Seconds** |
| **(19) Authors observe a deuteron beam polarization lifetime near 1000 s in (eb) the horizontal plane of (e) a magnetic storage ring (COSY).** |
| **(20) This long spin coherence time is maintained through (e&eb) a combination of beam bunching, electron cooling, (e&eb) sextupole field corrections, and (e&eb) the suppression of collective effects through (e&eb) beam current limits.** |
| **(21) This record lifetime is required for (e) a storage ring search for (e) an intrinsic electric dipole moment on (e&eb) the deuteron at a statistical sensitivity level approaching 10−29 e cm.** |
| **Synopsis: Deuterons Spin together for a Thousand Seconds July 28, 2016 Researchers set a new record for the in-plane spin-alignment lifetime of deuterons circulating in a magnetic storage ring. Synopsis figure E. Stephenson/Indiana University Center for Spacetime Symmetries Sime sentences and figures are deleted due to spatial restrictions.** |
| **(22) The electric dipole moment (EDM) of (e) a fundamental particle or a simple nucleus is believed to be (=) extremely small.** |
| **(23) One might therefore think that it isn't worth measuring this quantity. But it is. A nonzero EDM would signal (eb) a broken symmetry of (e) nature known as CP violation, which is required to (e) explain why the Universe has (e) more matter than antimatter.** |
| **(24) Edward Stephenson from Indiana University, Bloomington, and colleagues has now made a significant stride towards determining (eb) the EDM of the deuteron using (e) a deuteron beam circulating in (e&eb) a magnetic storage ring.** |

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| **(25) They have been able to preserve the alignment of the deuterons' spins for (e) a chunk of time that may be long enough to determine (eb) the deuteron's EDM.** |
| **(26) In experiments of this kind, the EDM can be determined by (e) measuring the rotation of (e) the initially aligned spins of (e) the beam particles as they travel around (e&eb) the ring.** |
| **(27) To obtain a measurement sufficiently accurate to explore (e&eb) CP violation predicted by (e) theories beyond the standard model, the spin alignment has to be (=) preserved for about 1000 s.** |
| **(28) But changes in particle speed—associated with (e&eb) position or direction variations among (e&eb) the particles in the beam—lead to (eb) a rapid loss of the alignment in the plane of the ring.** |
| **(29) Stephenson and co-workers have managed to retain (eb)the alignment of deuterons for the required 1000 s, which is 3 orders of magnitude longer (eb) than the previous record (obtained for an electron beam).** |
| **(30) Their success resulted from (e) capturing the beam into a tight bunch to reduce (e) position and direction spread and from (e) shaping the ring's magnetic field to(e) ensure all of the deuteron orbits around (e&eb) the ring have the same length, despite (e) direction variations. This research is published in Physical Review Letters. –Ana Lopes Ana Lopes is a Senior Editor of Physics.** |
| **Cotunneling Drag Effect in Coulomb-Coupled Quantum Dots A. J. Keller, J. S. Lim, David Sánchez, Rosa López, S. Amasha, J. A. Katine, Hadas Shtrikman, and D. Goldhaber-Gordon Phys. Rev. Lett. 117, 066602 – Published 4 August 2016 Physics logo See Synopsis: Coulomb Drag in a Double Dot** |
| **(31) In Coulomb drag, a current flowing in one conductor can induce (eb) a voltage across an adjacent conductor via (e&eb) the Coulomb interaction.** |
| **(32) The mechanisms yielding drag effects are not always understood, even though drag effects are (=) sufficiently general to be (=) seen in many low-dimensional systems.** |
| **(33) In this Letter, authors observe Coulomb drag in (eb) a Coulomb-coupled double quantum dot and, through (e&eb) both experimental and theoretical arguments, identify cotunneling as (=) essential to obtaining a correct qualitative understanding of the drag behavior.** |
| **Synopsis: Coulomb Drag in a Double Dot August 4, 2016 Electric current passing through a quantum dot can generate current in a nearby dot through a coordinated tunneling mechanism. Synopsis figureA. J. Keller et al., Phys. Rev. Lett. (2016)** |
| **(34) Electrons moving through (e&eb) a conductor can induce (eb) the movement of electrons of a nearby, but isolated, conductor.** |
| **(35) This so-called Coulomb drag is well understood in 2D layered materials, but the theory for nanowires and quantum dots does not (e) fully explain experiments.** |
| **(36) A new study of Coulomb drag in two nearby quantum dots shows (eb) that simultaneous tunneling of electrons—called cotunneling—can help explain the low-temperature behavior of (e) the system.** |

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| **(37) Electric current flows through (e&eb) a quantum dot by the tunneling of (e) electrons.** |
| **(38) When two dots are placed next to each other, the active, or lead, dot can drag current through (e&eb) the passive dot by (e) a process called sequential tunneling.** |
| **(39) In this case, an electron first tunnels onto (e&eb) the active dot, increasing (eb) the probability of tunneling onto (e&eb) the passive dot.** |
| **(40) Thus the passive dot follows (e&eb) the active dot's lead, but always one step behind.** |
| **(41) Sequential tunneling explains (eb) Coulomb drag observations at (eb) high temperature, but no studies have investigated cold temperatures where (e) quantum coherence effects may play a role.** |
| **(42) Andrew Keller from the California Institute of Technology, Pasadena, and his colleagues explored Coulomb drag at (eb) sub-Kelvin temperatures in a pair of quantum dots coupled by (e) capacitance.** |
| **(43) The team applied a current-driving voltage across (e&eb) one dot and detected a drag current in (eb) the other dot, even in regimes where (e) the driving voltage was very small.** |
| **(44) This conflicts with (e&eb) the sequential tunneling model, which predicts (eb) a threshold driving voltage below which (e) the drag current should be zero.** |
| **(45) The team showed that the inclusion of cotunneling into (eb) the drag theory explains the observations.** |
| **(46) Cotunneling could be important to efforts aimed at harvesting a unidirectional current from (e) a fluctuating voltage supply. This research is published in Physical Review Letters. –Michael Schirber Michael Schirber is a Corresponding Editor for Physics based in Lyon, France.** |
| **Maybe it wasn't the Higgs particle after all November 7, 2014 Maybe it wasn't the Higgs particle after all Associate Professor Mads Toudal Frandsen, University of Southern Denmark Credit: University of Southern Denmark** |
| **(47) Last year CERN announced the finding of a new elementary particle, the Higgs particle. But maybe it wasn't the Higgs particle, maybe it just looks like it. And maybe it is not alone. Many calculations indicate that the particle discovered last year in the CERN particle accelerator was (=) indeed the famous Higgs particle.** |
| **(48) Physicists agree that the CERN experiments did find a new particle that had never been seen before, but according to (e&eb) an international research team, there is (=) no conclusive evidence that the particle was indeed the Higgs particle.The research team has scrutinized the existing scientific data from CERN about the newfound particle and published their analysis in the journal Physical Review D. A member of this team is Mads Toudal Frandsen, associate professor at the Center for Cosmology and Particle Physics Phenomenology, Department of Physics, Chemistry and Pharmacy at the University of Southern Denmark. "The CERN data is generally taken as (=) evidence that the particle is the Higgs particle.** |

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| **(49) It is true that the Higgs particle can explain (eb) the data but there can be other explanations, we would also get this data from other particles", Mads Toudal Frandsen explains.** |
| **(50) The researchers' analysis does not (e) debunk the possibility that CERN has discovered the Higgs particle. That is still possible - but it is equally possible that it is a different kind of particle. "The current data is not precise enough to determine exactly what the particle is. It could be a number of other known particles", says Mads Toudal Frandsen. But if it wasn't the Higgs particle, which was found in CERN's particle accelerator, then what was it? "We believe that it may be a so-called techni-higgs particle. This particle is in some ways similar to the Higgs particle - hence half of the name", says Mads Toudal Frandsen.** |
| **(51) Although the techni-higgs particle and Higgs particle can easily be confused in experiments, they are two very different particles belonging to (e) two very different theories of how the universe was created.** |
| **(52) The Higgs particle is the missing piece in (eb) the theory called the Standard Model.** |
| **(53) This theory describes (eb) three of the four forces of nature. But it does not (e) explain what dark matter is - the substance that makes up (eb) most of the universe. A techni-higgs particle, if it exists, is (=) a completely different thing:** |
| **(54) "A techni-higgs particle is not (e) an elementary particle.** |
| **(55) Instead, it consists of so-called techni-quarks, which we believe are elementary. Techni-quarks may bind together in various ways to form (eb) for instance techni-higgs particles, while other combinations may form (eb) dark matter. Authors therefore expect to find several different particles at the LHC, all built by (e) techni-quarks", says Mads Toudal Frandsen.** |
| **(56) If techni-quarks exist, (eb) there must be a force to bind them together so that (e) they can form particles.** |
| **(57) None of the four known forces of nature (gravity, the electromagnetic force, the weak nuclear force and the strong nuclear force) are (=) any good at binding techni-quarks together.** |
| **(58) There must therefore be (=) a yet undiscovered force of nature. This force is called the the Technicolor force.** |
| **(59) What was found last year in CERN's accelerator could thus be (=) either the Higgs particle of the Standard Model or a light techni-higgs particle, composed of (e) two techni-quarks.** |
| **(60) Mads Toudal Frandsen believes that more data from CERN will probably be able to determine (eb) if it was a Higgs or a techni-higgs particle. If CERN gets an even more powerful accelerator, it will in principle be able to observe techni-quarks directly. More information: Technicolor Higgs boson in the light of LHC data. Phys. Rev. D 90, 035012th Alexander Belyaev, Matthew S. Brown, Roshan Foadi, and Mads T. Frandsen. journals.aps.org/prd/abstract/10.1103/PhysRevD.90.035012. On Arxiv: arxiv.org/abs/1309.2097 Journal reference: Physical Review D search and more info website: Provided by: University of Southern Denmark search and more info website** |

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| **The inflationary brane-antibrane universeClifford P. Burgess1,2, Mahbub Majumdar3, Detlef Nolte1, Fernando Quevedo3, Govindan Rajesh1 and Ren-Jie Zhang1 Published 13 August 2001 • Journal of High Energy Physics, Volume 2001, JHEP07(2001)** |
| **(61) Authors show how the motion through (e&eb) the extra dimensions of a gas of branes and antibranes can, under certain circumstances, produce (eb) an era of inflation as seen by (e) observers trapped on (e&eb) a 3-brane, with (e&eb) the inflaton being the inter-brane separation.** |
| **(62) Although most of discussion refers to (e) arbitrary ps, when we need to be specific they assume that they are (=) D-branes of type-II or type-I string theory.** |
| **(63) For realistic brane couplings, such as those arising in (eb) string theory, the inter-brane potentials are (=) too steep to inflate (eb) the universe for acceptably long times.** |
| **(64) However, for special regions of the parameter space of (e) brane-antibrane positions the brane motion is (=) slow enough for there to be (=) sufficient inflation.** |
| **(65) Inflation would be more generic in models where (e) the inter-brane interactions are much weaker.** |
| **(66) The spectrum of primordial density fluctuations predicted has (e) index n slightly less than 1, and acceptable amplitude, provided (e) that the extra dimensions have (e) linear size 1/r ~ 1012 GeV.** |
| **(67) Reheating occurs as in hybrid inflation, with (e&eb) the tachyonic instability of (e) the brane-antibrane system taking over for (e) small separations.** |
| **(68) The tachyon field can induce (eb) a cascade mechanism within which higher-dimension branes annihilate (e) into lower-dimension ones.** |
| **(69) They argue that such a cascade naturally stops with (e&eb) the production of 3-branes in (eb) 10-dimensional string theory.** |
| **Effect of the neutron lifetime on processes in the early universe A. V. Chechkin, A. V. Ivanchik, A. P. Serebrov, S. V. Bobashev First Online: 27 July 2016 DOI: 10.1134/S1063784216070069 Chechkin, A.V., Ivanchik, A.V., Serebrov, A.P. et al. Tech. Phys. (2016) 61: 1101. Doi: 10.1134/S1063784216070069** |
| **(70) The influence of the neutron lifetime on (e&eb) the abundance of light elements produced during (e&eb) the primordial nucleosynthesis early in (eb) the birth of the Universe is considered.** |
| **(71) Among light elements, namely, D, 3He, 4He, and 7Li, 4He proves to be (=) most sensitive to neutron lifetime τn.** |
| **(72) Astronomic data on the light element abundance also provide (eb) the best accuracy for 4He.** |
| **(73) The solution of a number of problems discussed in this paper requires (e) improving the accuracy of observations for the 4He abundance and refining the value of τn.** |

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| **Application of a Theory and Simulation based Convective Boundary Mixing model for AGB Star Evolution and Nucleosynthesis U. Battino, M. Pignatari, C. Ritter, F. Herwig, P. Denisenkov, J.W. Den Hartogh, R. Trappitsch, R. Hirschi, B. Freytag, F. Thielemann, B. Paxton (Submitted on 19 May 2016)** |
| **(74) The s-process nucleosynthesis in Asymptotic Giant Branch (AGB) stars depends on (e&eb) the modeling of (e) convective boundaries.** |
| **(75) Authors' present models and s-process simulations that adopt (e&eb) a treatment of convective boundaries based on (e) the results of hydrodynamic simulations and on (e&eb) the theory of mixing due to (e) gravity waves in (eb) the vicinity of convective boundaries.** |
| **(76) Hydrodynamics simulations suggest (eb) the presence of convective boundary mixing (CBM) at (eb) the bottom of the thermal pulse-driven convective zone.** |
| **(77) Similarly, convection-induced mixing processes are proposed for (e) the mixing below the convective envelope during (e&eb) third dredge-up where (e) the 13C pocket for (e) the s process in AGB stars forms.** |
| **(78) In this work apply apply a CBM model motivated by (e) simulations and theory to (e&eb) models with initial mass M = 2 and M = 3M?, and with (e&eb) initial metal content Z = 0:01 and Z = 0:02.** |
| **(79) As reported previously, the He-intershell abundance of 12C and 16O are increased by (e) CBM at the bottom of (e) pulse-driven convection zone.** |
| **(80) This mixing is affecting (e&eb) the 22Ne (alpha, n) 25Mg activation and (e&eb) the s-process effciency in the 13C-pocket.** |
| **(81) In our model CBM at the bottom of the convective envelope during (e&eb) the third dredgeup represents (eb) gravity wave mixing.** |
| **(82) Authors take further into account that hydrodynamic simulations indicate (eb) a declining mixing efficiency already about (e&eb) a pressure scale height from (e) the convective boundaries, compared to (e&eb) mixing-length theory.** |
| **(83) The final s-process abundances are characterized by (e) 0:36 < [s=Fe] < 0:78 and the heavy-to-light s-process ratio is (=) 0:23 < [hs=ls] < 0:45.** |
| **(84) Finally, authors compare results with (e&eb) stellar observations, pre-solar grain measurements and previous work. Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI:10.3847/0004-637X/827/1/30 Cite as: arXiv:1605.06159 [astro-ph.SR] (or arXiv:1605.06159v1 [astro-ph.SR] for this version)** |
| **MESA and NuGrid simulations of classical novae: CO and ONe nova nucleosynthesis Pavel A. Denissenkov (UVic), James W. Truran (Chicago), Marco Pignatari (Basel), Reto Trappitsch (Chicago), Christian Ritter (Frankfurt), Falk Herwig (UVic), Umberto Battino (Basel), Kiana Setoodehnia (McMaster), B. Paxton (UC Santa Barbara)** |

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| **(85) Classical novae are the result of (e) thermonuclear flashes of hydrogen accreted by (e) CO or ONe white dwarfs, leading (eb) eventually to the dynamic ejection of (e) the surface layers.** |
| **(86) These are observationally known to be (=) enriched in heavy elements, such as (=) C, O and Ne that must originate in (eb) layers below the H-flash convection zone.** |
| **(87) Building on their previous work, they now present stellar evolution simulations of (e) ONe novae and provide (eb) a comprehensive comparison of models with (e&eb) published ones.** |
| **(88) Some of authors models include exponential convective boundary mixing to account for (e) the observed enrichment of (e) the nova ejecta even when accreted material has (e) a solar abundance distribution.** |
| **(89) Models produce (eb) maximum temperature evolution profiles and nucleosynthesis yields (eb) in good agreement with (e&eb) models that generate (eb) enriched ejecta by assuming (eb) that the accreted material was pre-mixed.** |
| **(90) Authors confirm for (e) ONe novae the result we reported previously, i.e.\ we found that 3He could be produced {\it in situ} in (eb) solar-composition envelopes accreted with (e&eb) slow rates (M˙<10−10M⊙/yr) by (e) cold (TWD<107 K) CO WDs, and that convection was triggered by (eb) 3He burning before (e) the nova outburst in that case.** |
| **(91) In addition, authors now find that the interplay between the 3He production and (e&eb) destruction in the solar-composition envelope accreted with (e&eb) an intermediate rate, e.g.\ M˙=10−10M⊙/yr, by the 1.15M⊙ ONe WD with (e&eb) a relatively high initial central temperature, e.g.\ TWD=15×106 K, leads to (eb) the formation of a thick radiative buffer zone that separates (e&eb) the bottom of the convective envelope from (e) the WD surface. (Abridged) Subjects: Solar and Stellar Astrophysics (astro-ph.SR) DOI: 10.1093/mnras/stu1000 Cite as: arXiv:1303.6265 [astro-ph.SR] (or arXiv:1303.6265v2 [astro-ph.SR] for this version)** |
| **Hybrid C-O-Ne White Dwarfs as Progenitors of Diverse SNe Ia Pavel Denissenkov (UVic), James Truran (U. Chicago), Falk Herwig (UVic), Sam Jones (UVic), Bill Paxton (UC Santa Barbara), Ken'ichi Nomoto (U. Tokyo), Toshio Suzuki (Nihon U.), Hiroshi Toki (Osaka U.)(Submitted on 6 Nov 2014)** |
| **(92) When carbon is ignited off-center in a CO core of a super-AGB star (eb) its burning in a convective shell tends to propagate to (e&eb) the center.** |
| **(93) Whether the C flame will actually be able to reach (eb) the center depends on (e&eb) the efficiency of (e) extra mixing beneath (e&eb) the C convective shell.** |
| **(94) Whereas thermohaline mixing is too inefficient to interfere with the C-flame propagation, (eb) convective boundary mixing can prevent (e) the C burning from reaching (eb) the center.** |
| **(95) As a result, a C-O-Ne white dwarf (WD) is (=) formed, after the star has lost (e) its envelope.** |
| **(96) Such a "hybrid" WD has (e) a small CO core surrounded by (e) a thick ONe zone.** |

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| **(97) In the proposed 1D stellar evolution computations the hybrid WD is allowed to (eb) accrete C-rich material, as if it were in a close binary system and (e&eb) accreted H-rich material from (e) its companion with (e&eb) a sufficiently high rate at which the accreted H would be processed into (e&eb) He under (e&eb) stationary conditions, assuming (eb) that He could then be transformed into (e&eb) C.** |
| **(98) When the mass of the accreting WD approaches (eb) the Chandrasekhar limit, they find a series of convective Urca shell flashes associated with (e&eb) high abundances of 23Na and 25Mg.** |
| **(99) They are followed by off-center C ignition leading to (eb) convection that occupies (e) almost the entire star.** |
| **(100) To model the Urca processes, authors use the most recent well-resolved data for (e) their reaction and (e&eb) neutrino-energy loss rates.** |
| **(101) Because of the emphasized uncertainty of the convective Urca process in hybrid WD models of (e) SN Ia progenitors, authors consider a number of their potentially possible alternative instances for (e) different mixing assumptions, all of which reach (eb) a phase of explosive C ignition, either off or in (eb) the center.** |
| **(102) Our hybrid SN Ia progenitor models have (e) much lower C to O abundance ratios at (eb) the moment of the explosive C ignition than (e) their pure CO counterparts, which may explain (eb) the observed diversity of the SNe Ia. subjects: Solar and Stellar Astrophysics (astro-ph.SR)cite as: arXiv:1411.1471 [astro-ph.SR](or arXiv:1411.1471v1 [astro-ph.SR] for this version)** |
| **Rejecting proposed dense-matter equations of state with quiescent low-mass X-ray binaries Sebastien Guillot, Robert E. Rutledge (Submitted on 15 Sep 2014 (v1), last revised 20 Oct 2014 (this version, v2))** |
| **(103) Neutrons stars are unique laboratories to (e) discriminate between the various proposed equations of (e) state of matter at and above nuclear density.** |
| **(104) One sub-class of neutron stars - those inside (eb) quiescent low-mass X-ray binaries (qLMXBs) - produce (eb) a thermal surface emission from (e) which the neutron star radius (R_NS) can be measured, using (e) the widely accepted observational scenario for qLMXBs, assuming (eb) unmagnetized H atmospheres.** |
| **(105) In a combined spectral analysis, this work first reproduces (eb) a previously published measurement of the \rns, assumed to be (=) the same for all neutron stars, using (e) a slightly expanded data set.** |
| **(106) The radius measured is R_NS = 9.4 +/-1.2 km. On the basis of spectral analysis alone, this measured value is not affected by imposing (e&eb) an assumption of causality in the core.** |
| **(107) However, the assumptions underlying this R_NS measurement would be falsified by (e) the observation of any neutron star with (e&eb) a mass >2.6 Msun, since radii <11 km would be rejected if (e) causality is assumed, which would exclude (eb) most of the R_NS parameter space obtained in this analysis.** |

**(108) Finally, this work directly tests (e&eb) a selection of dense matter equations of states: WFF1, AP4, MPA1, PAL1, MS0, and three versions of equations of state produced through (e&eb) chiral effective theory.**

**(109) Two of those, MSO and PAL1, are rejected at (eb) the 99% confidence level, accounting for (e) all quantifiable uncertainties, while (e) the other cannot be excluded at >99% certainty. Subjects: High Energy Astrophysical Phenomena (astro-ph.HE) DOI: 10.1088/2041-8205/796/1/L3 Cite as: arXiv:1409.4306 [astro-ph.HE] (or arXiv:1409.4306v2 [astro-ph.HE] for this version**

## SECTION NINE

## Analytical Computation Of The Mercury Perihelion Precession Via The Relativistic Gravitational Law And Comparison With General Relativity

### INTRODUCTION—VARIABLES USED

**Analytical computation of the Mercury perihelion precession via the relativistic gravitational law and comparison with general relativity Athanassios S. Fokas, Constantinos G. Vayenas, D. Grigoriou**

(1) Let r (φ) denote (=) the orbit of Mercury.
(2) Authors compare the formulae obtained via (e&eb) general relativity for (e) r(φ) and for the corresponding (e&eb) perihelion precession angle Δφ, with (e&eb) the formulae obtained via (e&eb) the relativistic gravitational law, F=GMmγ6/r2.
(3) The latter law can be derived from (e) Newton's gravitational law by employing (e) the gravitational rather than the rest masses of (e) the Sun and Mercury.
(4) Remarkably, it is found that the two expressions for (e) the dominant part of r (φ) are identical and (e&eb) that the two expressions for (e) Δφ are also identical. Subjects: General Physics (physics.gen-ph) Cite as: arXiv: 1509.03326 [physics.gen-ph] (or arXiv: 1509.03326v2 [physics.gen-ph] for this version)

**A Bohr-type model with gravity as the attractive force C.G. Vayenas, S. Souentie, A. Fokas**

(5) Authors formulate a Bohr-type rotating particle model for (e) three light particles of the same rest mass, forming (eb) a bound rotational state under the influence of (e&eb) their gravitational attraction, in the same way that electrostatic attraction leads to (eb) the formation of a bound proton-electron state in (eb) the classical Bohr model of the H atom.
(6) By using special relativity, the equivalence principle and the de Broglie wavelength equation, authors find that when (e) the three rotating particles have (e) the rest masses of neutrinos or antineutrinos then (eb) surprisingly the rest mass of the rotating state has (e) the rest mass of the stable baryons, i.e. of the proton and the neutron.
(7) This rest mass is due almost exclusively to (e) the kinetic energy of the rotating neutrinos.
(8) The results are found to be consistent with (e&eb) the theory of general relativity.
(9) Predictions for the properties of these bound rotational states are compared with (e&eb) experimental values. Subjects:General Physics (physics.gen-ph) Journal reference: Physica A, 405 (2014) 360-379 DOI:10.1016/j.physa.2014.03.045 Cite as: arXiv:1306.5979

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| [physics.gen-ph] (or arXiv:1306.5979v4 [physics.gen-ph] for this version) | |
| **NOTATION** | |
| **Module One**<br><br>Let r (φ) denote (=) the orbit of Mercury | |
| $G_{13}$ : Category one of r (φ)<br><br>$G_{14}$ : Category two of SAS<br><br>$G_{15}$ : Category three of SAS | |
| $T_{13}$ : Category one of orbit of Mercury<br><br>$T_{14}$ : Category two of SAS<br><br>$T_{15}$ : Category three of SAS | |
| **Module Two**<br><br>Authors compare the formulae obtained via general relativity fo r(φ) and for the corresponding perihelion precession angle Δφ, with (e&eb) the formulae obtained via (e&eb) the relativistic gravitational law, F=GMmγ6/r2 | |
| $G_{16}$ : Category one of **formulae obtained via general relativity fo r(φ) and for the corresponding perihelion precession angle Δφ;** formulae obtained via (e&eb) the relativistic gravitational law, F=GMmγ6/r2<br><br>$G_{17}$ : Category two of SAS<br><br>$G_{18}$ : Category three of SAS | |
| $T_{16}$ : Category one of formulae obtained via (e&eb) the relativistic gravitational law, F=GMmγ6/r2; **formulae obtained via general relativity fo r(φ) and for the corresponding perihelion precession angle Δφ**<br><br>$T_{17}$ : Category two of SAS<br><br>$T_{18}$ : Category three of SAS | |
| **Module three**<br><br>Authors compare the formulae obtained via general relativity fo r(φ) and for the corresponding perihelion precession angle Δφ, with (e&eb) the formulae obtained via (e&eb) the relativistic gravitational law, F=GMmγ6/r2 | |
| $G_{20}$ : Category one of **formulae obtained;** general relativity fo r(φ) and for the corresponding perihelion precession angle Δφ, | |

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| $G_{21}$ : Category two of SAS<br><br>$G_{22}$ : Category three of SAS | |
| $T_{20}$ : Category one of general relativity fo r(φ) and for the corresponding perihelion precession angle Δφ,;**formulae obtained**<br><br>$T_{21}$ : Category two of SAS<br><br>$T_{22}$ : Category three of SAS | |
| **Module four**<br><br>Authors compare the formulae obtained via general relativity fo r(φ) and for the corresponding perihelion precession angle Δφ, with (e&eb) the formulae obtained via (e&eb) the relativistic gravitational law, F=GMmγ6/r2 | |
| $G_{24}$ : Category one of **formulae;** relativistic gravitational law, F=GMmγ6/r2<br><br>$G_{25}$ : Category two of SAS<br><br>$G_{26}$ : Category three of SAS | |
| $T_{24}$ : Category one of relativistic gravitational law, F=GMmγ6/r2; **formulae**<br><br>$T_{25}$ : Category two of SAS<br><br>$T_{26}$ : Category three of SAS | |
| **Module five**<br><br>The latter law can be derived from (e) Newton's gravitational law by employing (e) the gravitational rather than the rest masses of (e) the Sun and Mercury | |
| $G_{28}$ : Category one of Newton's gravitational law by employing (e) the gravitational rather than the rest masses of (e) the Sun and Mercury<br><br>$G_{29}$ : Category two of SAS<br><br>$G_{30}$ : Category three of SAS | |
| $T_{28}$ : Category one of latter law can be derived<br><br>$T_{29}$ : Category two of SAS<br><br>$T_{30}$ : Category three of SAS | |
| **Module six**<br><br>The latter law can be derived from Newton's gravitational law by employing (e) the gravitational rather than the rest masses of (e) the Sun and Mercury | |

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| $G_{32}$ : Category one of gravitational rather than the rest masses of (e) the Sun and Mercury<br><br>$G_{33}$ : Category two of SAS<br><br>$G_{34}$ : Category three of SAS | |
| $T_{32}$ : Category one of latter law can be derived from Newton's gravitational law by employing<br><br>$T_{33}$ : Category two of SAS<br><br>$T_{34}$ : Category three of SAS | |
| **Module seven**<br><br>The latter law can be derived from Newton's gravitational law by employing the gravitational rather than the rest masses of (e) the Sun and Mercury | |
| $G_{36}$ : Category one of Sun and Mercury (use average for calculation take all the three categories as equivalent)<br><br>$G_{37}$ : Category two of SAS<br><br>$G_{38}$ : Category three of SAS | |
| $T_{36}$ : Category one of latter law can be derived from Newton's gravitational law by employing the gravitational rather than the rest masses<br><br>$T_{37}$ : Category two of SAS<br><br>$T_{38}$ : Category three of SAS | |
| **Module eight**<br><br>Remarkably, it is found that the two expressions for (e) the dominant part of r (φ) are identical and (e&eb) that the two expressions for (e) Δφ is also identical.<br><br>Subjects: General Physics (physics.gen-ph) Cite as: arXiv: 1509.03326 [physics.gen-ph] (or arXiv: 1509.03326v2 [physics.gen-ph] for this version) | |
| $G_{40}$ : Category one of dominant part of r (φ) is identical and (e&eb) that the two expressions for (e) Δφ is also identical.<br><br>$G_{41}$ : Category two of SAS<br><br>$G_{42}$ : Category three of SAS | |
| $T_{40}$ : Category one of two expressions<br><br>$T_{41}$ : Category two of SAS<br><br>$T_{42}$ : Category three of SAS | |

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| **Module Nine**<br><br>Remarkably, it is found that the two expressions for the dominant part of r (φ) are identical and (e&eb) that the two expressions for (e) Δφ is also identical | |
| $G_{44}$ : Category one of **two expressions for the dominant part of r (φ) are identical;** two expressions for Δφ is also identical<br><br>$G_{45}$ : Category two of SAS<br><br>$G_{46}$ : Category three of SAS | |
| $T_{44}$ : Category one of two expressions for Δφ is also identical; **two expressions for the dominant part of r (φ) are identical**<br><br>$T_{45}$ : Category two of SAS<br><br>$T_{46}$ : Category three of SAS | |

## SECTION TEN

## A Bohr-Type Model With Gravity As The Attractive Force

**INTRODUCTION—VARIABLES USED**

**A Bohr-type model with gravity as the attractive force C.G. Vayenas, S. Souentie, A. Fokas**

(1) Authors formulate a Bohr-type rotating particle model for (e) three light particles of the same rest mass, forming (eb) a bound rotational state under the influence of (e&eb) their gravitational attraction, in the same way that electrostatic attraction leads to (eb) the formation of a bound proton-electron state in (eb) the classical Bohr model of the H atom.
(2) By using special relativity, the equivalence principle and the de Broglie wavelength equation, authors find that when (e) the three rotating particles have (e) the rest masses of neutrinos or antineutrinos then (eb) surprisingly the rest mass of the rotating state has (e) the rest mass of the stable baryons, i.e. of the proton and the neutron.
(3) This rest mass is due almost exclusively to (e) the kinetic energy of the rotating neutrinos.
(4) The results are found to be consistent with (e&eb) the theory of general relativity.
(5) Predictions for the properties of these bound rotational states are compared with (e&eb) experimental values. Subjects:General Physics (physics.gen-ph) Journal reference: Physica A, 405 (2014) 360-379 DOI:10.1016/j.physa.2014.03.045 Cite as: arXiv:1306.5979 [physics.gen-ph] (or arXiv:1306.5979v4 [physics.gen-ph] for this version)

**NOTATION**

**Module One**

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| Remarkably, it is found that the two expressions for (e) the dominant part of r (φ) are identical and (e&eb) that the two expressions for (e) Δφ is also identical.<br><br>Subjects: General Physics (physics.gen-ph) Cite as: arXiv: 1509.03326 [physics.gen-ph] (or arXiv: 1509.03326v2 [physics.gen-ph] for this version) | |
| $G_{13}$ : Category one of dominant part of r (φ) is identical and (e&eb) that the two expressions for (e) Δφ is also identical.<br><br>$G_{14}$ : Category two of SAS<br><br>$G_{15}$ : Category three of SAS | |
| $T_{13}$ : Category one of two expressions<br><br>$T_{14}$ : Category two of SAS<br><br>$T_{15}$ : Category three of SAS | |
| **Module Two**<br><br>Remarkably, it is found that the two expressions for the dominant part of r (φ) are identical and (e&eb) that the two expressions for (e) Δφ is also identical.<br><br>Subjects: General Physics (physics.gen-ph) Cite as: arXiv: 1509.03326 [physics.gen-ph] (or arXiv: 1509.03326v2 [physics.gen-ph] for this version | |
| $G_{16}$ : Category one of **two expressions for the dominant part of r (φ) are identical;** two expressions for (e) Δφ is also identical<br><br>$G_{17}$ : Category two of SAS<br><br>$G_{18}$ : Category three of SAS | |
| $T_{16}$ : Category one of two expressions for (e) Δφ is also identical **;two expressions for the dominant part of r (φ) are identical**<br><br>$T_{17}$ : Category two of SAS<br><br>$T_{18}$ : Category three of SAS | |
| **Module three**<br><br>Authors formulate a Bohr-type rotating particle model for (e) three light particles of the same rest mass, forming (eb) a bound rotational state under the influence of (e&eb) their gravitational attraction, in the same way that electrostatic attraction leads to (eb) the formation of a bound proton-electron state in (eb) the classical Bohr model of the H atom | |
| $G_{20}$ : Category one ofthree light particles of the same rest mass, forming (eb) a bound rotational state under the influence of (e&eb) their gravitational attraction, in the same way that electrostatic attraction leads to | |

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| (eb) the formation of a bound proton-electron state in (eb) the classical Bohr model of the H atom<br><br>$G_{21}$ : Category two of SAS<br><br>$G_{22}$ : Category three of SAS | |
| $T_{20}$ : Category one of Bohr-type rotating particle model<br><br>$T_{21}$ : Category two of SAS<br><br>$T_{22}$ : Category three of SAS | |
| **Module four**<br><br>Authors formulate a Bohr-type rotating particle model for three light particles of the same rest mass, forming (eb) a bound rotational state under the influence of (e&eb) their gravitational attraction, in the same way that electrostatic attraction leads to (eb) the formation of a bound proton-electron state in (eb) the classical Bohr model of the H atom | |
| $G_{24}$ : Category one of formulate a Bohr-type rotating particle model for three light particles of the same rest mass<br><br>$G_{25}$ : Category two of SAS<br><br>$G_{26}$ : Category three of SAS | |
| $T_{24}$ : Category one ofbound rotational state under the influence of (e&eb) their gravitational attraction, in the same way that electrostatic attraction leads to (eb) the formation of a bound proton-electron state in (eb) the classical Bohr model of the H atom<br><br>$T_{25}$ : Category two of SAS<br><br>$T_{26}$ : Category three of SAS | |
| **Module five**<br><br>Authors formulate a Bohr-type rotating particle model for three light particles of the same rest mass, forming a bound rotational state under the influence of (e&eb) their gravitational attraction, in the same way that electrostatic attraction leads to (eb) the formation of a bound proton-electron state in (eb) the classical Bohr model of the H atom | |
| $G_{28}$ : Category one of **Bohr-type rotating particle model for three light particles of the same rest mass, forming a bound rotational state;** their gravitational attraction, in the same way that electrostatic attraction leads to (eb) the formation of a bound proton-electron state in (eb) the classical Bohr model of the H atom<br><br>$G_{29}$ : Category two of SAS<br><br>$G_{30}$ : Category three of SAS | |

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| $T_{28}$ : Category one of their gravitational attraction, in the same way that electrostatic attraction leads to (eb) the formation of a bound proton-electron state in (eb) the classical Bohr model of the H atom **;Bohr-type rotating particle model for three light particles of the same rest mass, forming a bound rotational state**<br><br>$T_{29}$ : Category two of SAS<br><br>$T_{30}$ : Category three of SAS | |
| **Module six**<br><br>Authors formulate a Bohr-type rotating particle model for three light particles of the same rest mass, forming a bound rotational state under the influence of their gravitational attraction, in the same way that electrostatic attraction leads to (eb) the formation of a bound proton-electron state in (eb) the classical Bohr model of the H atom | |
| $G_{32}$ : Category one of Bohr-type rotating particle model for three light particles of the same rest mass, forming a bound rotational state under the influence of their gravitational attraction, in the same way that electrostatic attraction; **formation of a bound proton-electron state in (eb) the classical Bohr model of the H atom**<br><br>$G_{33}$ : Category two of SAS<br><br>$G_{34}$ : Category three of SAS | |
| $T_{32}$ : Category one of **formation of a bound proton-electron state in (eb) the classical Bohr model of the H atom;** Bohr-type rotating particle model for three light particles of the same rest mass, forming a bound rotational state under the influence of their gravitational attraction, in the same way that electrostatic attraction<br><br>$T_{33}$ : Category two of SAS<br><br>$T_{34}$ : Category three of SAS | |
| **Module seven**<br><br>Authors formulate a Bohr-type rotating particle model for three light particles of the same rest mass, forming a bound rotational state under the influence of their gravitational attraction, in the same way that electrostatic attraction leads to the formation of a bound proton-electron state in (eb) the classical Bohr model of the H atom | |
| $G_{36}$ : Category one of Bohr-type rotating particle model for three light particles of the same rest mass, forming a bound rotational state under the influence of their gravitational attraction, in the same way that electrostatic attraction leads to the formation of a bound proton-electron state<br><br>$G_{37}$ : Category two of SAS<br><br>$G_{38}$ : Category three of SAS | |

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| $T_{36}$ : Category one of classical Bohr model of the H atom<br><br>$T_{37}$ : Category two of SAS<br><br>$T_{38}$ : Category three of SAS | |
| **Module eight**<br><br>By using special relativity, the equivalence principle and the de Broglie wavelength equation, authors find that when (e) the three rotating particles have (e) the rest masses of neutrinos or antineutrinos then (eb) surprisingly the rest mass of the rotating state has (e) the rest mass of the stable baryons, i.e. of the proton and the neutron | |
| $G_{40}$ : Category one of special relativity, the equivalence principle and the de Broglie wavelength equation<br><br>$G_{41}$ : Category two of SAS<br><br>$G_{42}$ : Category three of SAS | |
| $T_{40}$ : Category one ofwhen (e) the three rotating particles have (e) the rest masses of neutrinos or antineutrinos then (eb) surprisingly the rest mass of the rotating state has (e) the rest mass of the stable baryons, i.e. of the proton and the neutron<br><br>$T_{41}$ : Category two of SAS<br><br>$T_{42}$ : Category three of SAS | |
| **Module Nine**<br><br>By using special relativity, the equivalence principle and the de Broglie wavelength equation, authors find that when (e) the three rotating particles have (e) the rest masses of neutrinos or antineutrinos then (eb) surprisingly the rest mass of the rotating state has (e) the rest mass of the stable baryons, i.e. of the proton and the neutron | |
| $G_{44}$ : Category one ofthree rotating particles have (e) the rest masses of neutrinos or antineutrinos then (eb) surprisingly the rest mass of the rotating state has (e) the rest mass of the stable baryons, i.e. of the proton and the neutron<br><br>$G_{45}$ : Category two of SAS<br><br>$G_{46}$ : Category three of SAS | |
| $T_{44}$ : Category one of special relativity, the equivalence principle and the de Broglie wavelength equation<br><br>$T_{45}$ : Category two of SAS<br><br>$T_{46}$ : Category three of SAS | |

References:

Given in abstracts